Cataluminescence in Er-Substituted Perovskites

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Thermophotovoltaic devices have promising applications for energy conversion. However, current conversion efficiency of chemical energy to light is very low, limited by the competing process of heat dissipation released as black body radiation. From a fundamental point of view, the direct conversion of chemical energy into light without this detour is possible. This so-called cataluminescence from methanol combustion over Er-substituted SrTiO$_3$ with high efficiency is demonstrated. The catalytically active quaternary perovskites Er$_{0.15}$La$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_3$ - δ exsolute and reabsorb metallic Cu particles onto the surface in reducing and oxidizing conditions, respectively. Thus, it is able to manipulate the surface structure and investigate its influence on the catalytic as well as luminescent properties. The fuel to air ratio around the stoichiometry point changes the conditions from reducing to oxidizing and thereby alters the surface properties. This is evidenced by post mortem X-ray diffraction and X-ray photoemission as well as operando optical spectroscopy. Cataluminescence takes place under oxidizing conditions (lean fuel to air mixture) on the Er-perovskite oxide with a strong selective near infrared emission, while reducing conditions stimulate formation of plasmonic Cu-nanoparticles, which emit black body radiation.

1. Introduction

Thermophotovoltaic converters transform chemical energy of basically any fuel to electricity. They are rather simple, and particularly suited to deliver electrical power to remote places for outdoor activities and military applications.[1] These devices traditionally consist of a burner, which combusts the fuel at high temperatures to generate light (chemical energy to light conversion), and a photovoltaic cell, which converts the light into electricity. The generated light is mainly black body radiation with main energy distribution in the far infrared region. As photovoltaic cells convert light only above the threshold energy, the efficiency of such a thermophotovoltaic device is very low. In the past, selective emitters were developed to increase the efficiency, but with limited success (thermoluminescence).[2] Circumventing heat emission (i.e., lower entropy production) may markedly increase the efficiency, extending the applicability of such emitters for energy conversion. Direct conversion of chemical energy to light is known in nature as chemiluminescence. Cataluminescence is a specific type of chemiluminescence that occurs on solid catalysts during heterogeneous fuel oxidation reactions. The phenomenon stems from combustion-formed radicals present at the surface of the light emitting material (in contrast to the outdated gas lights, where the radicals are formed in the gas phase).[3] It is the main non-thermal factor that determines light emission spectrum, first observed on TiO$_2$.[4] Since then, the effect was used to develop luminescence sensors for various chemical compounds,[5–7] which suggests another potential application, as an in situ probe of ongoing catalytic oxidation reactions. Understanding the phenomenon carries tremendous potential from the viewpoint of fundamental science and industry, which can be extrapolated from improving efficiency of methanol combustion to combustion of other fuels.

Thus, the first goal is to develop a material that fulfills the criteria of being catalytically active for methanol oxidation (Equation (1)) while simultaneously possessing electronic states that make use of the released reaction energy of about 1.1 eV by selective emission rather than heat generation.

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\text{CH}_3\text{OH} + \frac{1}{2} \text{O}_2 \rightarrow \text{CO}_2 + 2\text{H}_2\text{O}, \quad \Delta H = -740 \text{ kJ mol}^{-1} \tag{1}
\]
Typical methanol oxidation catalysts are based on precious metal nanoparticles deposited on oxide supports. In addition, rare earth oxides from the lanthanide series are well known for their thermoluminescence properties in the near infrared (NIR) region. For this purpose, we explored the combination of Ru, Pd, and Cu nanoparticles on rare earth supports in preliminary studies. However, we observed no or very inefficient selective emission and the typically low melting points of the precious metal nanoparticles (bulk Cu = 1084 °C) that resulted in a complete loss of catalytic activity already after a short reaction period (< minutes) assigned to sintering of the particles on the surface. 

Inspired by recent developments of oxide based catalysts, we propose a different pathway: metal oxide based perovskites with the general formula ABO$_3$ are well known for their flexibility in implementing foreign ions into the crystal lattice without major structural reorganization. Depending on their ionic radii, dopants can be introduced on the A- as well as on the B-sites leading to a general formula of AA′BB′O$_x$. If present, different redox potentials enable exsolution of these dopants as metal nanoparticles on the surface in a reductive atmosphere, with subsequent reabsorption into the crystal lattice in an oxidative atmosphere. This method differs from a simple post synthetic impregnation procedure and generates strong crystal interfaces, as evidenced by, for example, higher thermal stability of the nanoparticles on the surface.

In this paper we demonstrate the cataluminescence in Er- and Cu-substituted SrTiO$_3$ (Er$_{0.15}$La$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_{3–δ}$, Er-LaSTO). The paper is organized as follows: first section describes the general idea and first outcomes of cataluminescence from methanol combustion. To distinguish the physical effects of black body radiation, we make use of the peculiar properties of the catalytically active quaternary perovskites, which reversibly exsolve and absorb metallic Cu particles onto the surface when switching from oxidizing to reducing conditions. We characterize the material using a variety of operando and post-mortem methods. Particular emphasis is put on the exsolution process, which is followed by in situ plasmon resonance spectroscopy. During combustion, changing fuel to air ratio ("lambda ratio") around the stoichiometry point from reducing to oxidizing conditions triggers the formation of surface nanoparticles. With this, we are able to link the surface transformations to the emission transition (selective NIR-emission to black-body emission) due to the change of the lambda ratio. This is a strong indication that cataluminescence takes place on the Er-perovskite with a strong selective NIR-emission, while under reducing conditions black body radiation is emitted from catalytically active Cu nanoparticles.

2. Results

The catalytically active quaternary perovskite oxide Er$_{0.15}$La$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_{3–δ}$ (Er-LaSTO) was prepared and characterized as described in Section 5. Briefly, the material is a porous powder consisting of polycrystalline ErLaSTO in the perovskite and double perovskite structure (Figure 6) with a complex surface structure depending on the applied conditions. The powder is put into a modified Harrick cell for catalysis allowing operando UV–vis diffuse reflectance and NIR spectroscopy measurements (see Section 5 for details). The sample can be exposed to flows of different gases, such as air and hydrogen to study the redox behavior, and to air/methanol mixtures for catalytic combustion. The material catalyzes the methanol oxidation reaction (Equation (1)) over the whole experimentally accessible oxygen to methanol ratio $r = O_2/CH_3OH$ of 1.38 < $r$ < 4.6. As the reaction is highly exothermic, it is self-sustaining after ignition at the temperatures around 400 °C. It follows that the temperature on the catalyst is substantially higher. The infrared radiation emitted from the catalyst is used to estimate its temperature. Neglecting the emissivity, a black-body model is applied to the NIR spectra. An emission of similar absolute intensity is only obtained if a temperature around 700–800 °C is assumed (see Figure 1 and Section 5, Figure 5). However, the NIR emission spectra during catalysis from Er-substituted perovskite strongly differ from a black-body spectrum such as observed from the same sample if only heated (blue/red dots in Figure 1). Some features of the cataluminescence spectrum match that of the cathodo/thermo luminescence spectrum of Er$_2$O$_3$. With this, the spectral features can be assigned unambiguously to luminescent emission from Er-ions which also appear in the diffuse reflectance spectrum (Figure 1).

The selective emission is only present under catalytic conditions. External heating leads to black-body radiation only. To semi-quantify the effect, we tuned the emission intensity during catalysis (blue dots in Figure 1) by adjusting the fuel to air ratio in the mixture to the intensity of the thermal emission from the same material exposed to air only at 760 °C, which is the maximum temperature of the sample cell (red dots in Figure 1). The difference between the two measurements gives the efficiency of the emission enhancement in this spectral range by cataluminescence, which is of the order of 70%.

However, the strong enhancement alone is not sufficient enough to prove cataluminescence. The heat release confined to active sites might heat up the surface locally and induce thermoluminescence. The orange curve in Figure 1 is a simulation of black-body radiation at around 1200 °C divided by a factor of 20, that is, assuming only 5% of the surface is active. The slope of the curve at smaller wavelengths is reproduced, and the selective emission may be an effect from thermoluminescence.[2] To counter this argument and to provide evidence of a direct energy transport from the catalytically active site to the luminescent Er-ions, we make use of the peculiar redox behavior of the material. The B site in ABO$_3$ perovskites, here Ti, can be partially substituted by transition metals such as Cu without major changes to the crystal structure as well as electronic properties in oxidizing conditions.[11,14] It was found that the material exsolves metallic Cu particles onto the surface upon reduction at high temperatures. The same behavior follows upon hydrogen exposure at elevated temperature, as confirmed by temperature programmed reduction (TPR), post-mortem X-ray diffraction (XRD) and X-ray photoelectron spectroscopy (XPS) measurements on Er-substituted perovskites (Section 5).

The changes observed by XPS post catalysis indicate that also the conditions during catalysis restructure the surface of the Cu-doped catalysts (Figure 6). However, operando spectroscopy is required, as the surface studied post mortem will not resemble the active surface under catalytic conditions. Optical spectroscopy in the UV–vis range has yielded information on the formation of Cu-nanoparticles, and is rather easily implemented into the experimental setup (see Section 5). The transformation
Infrared radiation emitted from a solid catalyzing methanol combustion may originate from black-body radiation as a result of the heat of combustion and/or from cataluminescence, that is, the transfer of the chemical energy from combustion to luminescent centers. To distinguish the two phenomena, we measure the infrared radiation emitted purely by heating the sample, and during catalytic combustion (see Section 5 and Figure 5 for technical details). Left panel is the NIR-emission from cataluminescence (blue dots) and thermal emission at 760 °C (red dots) of Er_{0.15}La_{0.15}Sr_{0.55}Ti_{0.95}Cu_{0.05}O_{3-\delta}. The red and yellow lines are simulations of black-body radiation at 720 and 1200 °C, respectively. The right panel compares the black-body subtracted cataluminescence ($T_{BB}$ = 1473K) with the corresponding UV–vis–NIR diffuse reflectance spectrum. The broad negative peak around 600 nm is the spectral region of Cu-plasmon absorption. The sharp peaks are assigned to optical transitions in Er$^{3+}$ ions in accordance with Er$^{3+}$ in SrF$_2$,[15] and Er$^{3+}$ in sesquioxides.[2,16]

Operando UV–vis diffuse reflectance spectroscopy on Er$_{0.15}$La$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_{3-\delta}$ during hydrogen exposure from 20 to 650 °C. The right graph shows the subsequently recorded spectra after first exposing to pure air at 650 °C and to methanol / air mixtures at oxidizing ($r = 2.0$) and reducing ($r = 1.8$) conditions. The latter is performed at a cell temperature of 400 °C. The evolution of the peak at 585 nm is indicative of plasmon absorption in Cu-nanoparticles. The absolute absorbance spectrum of the hydrogenated sample measured ex situ is shown in Figure 1.

is evidenced by the color change of the samples from yellow to gray upon reduction, indicative of light absorption by metallic nanoparticles. The operando UV–vis diffuse reflectance spectroscopy (Figure 2) shows the evolution of the absorbance peak during heating of an as prepared ErLSTO sample in hydrogen atmosphere. In good agreement with TPR experiments and XPS analysis (see Figure 6), optical changes indicating hydrogen reacting with the material appear already at temperatures as low as 200 °C. However, the typical peak around 580 nm, indicative of plasmon excitation in Cu nanoparticles,[17–20] develops at higher temperatures. This coincides with the post mortem XRD analysis, detecting Cu-nanoparticles after hydrogen treatment at $T > 650$°C. The exsolution thus proceeds in two steps: emerging of surface copper atoms at low temperature, which is followed by diffusion and agglomeration to Cu-nanoparticles at higher temperatures. This process is assumed to be reversible. Exposing the freshly reduced sample to air at 900 K reduces the plasmon peak markedly, but not entirely. In particular, the shoulder at around 500 nm remains, which is assigned to CuO.[20] Full reversibility is expected to be reached at higher temperatures than possible in the optical cell.

The optical setup allows the measurement of absorbance changes under catalytic conditions. The intensity of the plasmon excitation depends markedly on the air-fuel ratio: the peak develops when changing from oxidizing ($r = 2.0$) to reducing ($r = 1.8$) conditions during methanol combustion, without any other spectral changes. This is a strong evidence for the formation/growth of nanoparticles under reducing conditions. We analyzed the products using a gas IR-spectrometer with particular focus on CO$_2$ and CO yields. The CO$_2$ yield is directly proportional to the rate of conversion. It is smaller under reducing conditions than under oxidizing conditions, that is, the conversion rate is slower. The CO yield is the opposite. As CO is a strong reducing agent, it can drive the exsolution of Cu and formation of Cu nanoparticles regardless of the overall oxidizing stoichiometry ratio.

The structural modifications coincide with the infrared emission, which changes around the ratio $1.8 < r < 2.0$ from selective emission of Er-ions to a spectrum similar to black-body radiation (Figure 3). The effect is reversible and solely dependent...
on the $r$ ratio. The methanol flow, that is, the effective heat release over the catalysts has little effect on the spectral shape as has the exact lambda ratio, as long as $r$ does not approach the critical value around 1.9. Cu-nanoparticles are known to be catalytically active, and are preferential as reaction sites under reducing conditions, as opposed to the sites in the vicinity of the Er-ions as is the case under oxidizing conditions. The heat evolved is then emitted mainly as black-body radiation.

3. Discussion

The results imply that the radiation released during catalytic combustion is markedly shifted toward more usable short wavelengths by pumping and radiative de-excitation of Er$^{3+}$ ions. The question arises about how the energy is transferred from the catalytic centers to the Er-ions during oxidizing conditions, and how the corresponding electrons are excited. The oxidation of methanol can be described as a series of redox-processes: six electrons are transferred from methanol to the oxygen atoms bound in CO$_2$ and water. Without energy transfer, as is the case on purely metallic catalysts, the corresponding electrons, originally at high chemical potential, dissipate their energy in the metal by a heating process. The same phenomenon was observed on the pre-experiments, where the combustion took place over catalytically active Cu. The origin of the selective emission is the radiative de-excitation of these Er$^{3+}$ ions. The averaged chemical potential difference per electron gained from methanol combustion is around 1.2 eV, which is just sufficient to pump the $^1I_{15/2}$ states. The nature of the pumping (energy transfer, ET) cannot be clarified in this paper. It is well known that the thermoluminescence efficiency depends very sensitively on the absolute position of the 4f electron energy levels with respect to the conduction and to the valence bands. This effect may be combined with the exciton formation and emission as proposed for the non-selective cataluminescence mechanism. That is, the excited radicals formed on catalytically active but non-plasmonic Cu create excitons, whose energy is transferred to the Er-ions. It is not expected that the energy transfer is perfect, that is, the catalytically active sites will heat up in any case. The observed temperature increase is then an unwanted side effect. Under reducing conditions, the copper particles become plasmonic and quench the hot electrons resulting in heat emission, that is, black body radiation with maximum at lower energies (Figure 4).

Further mechanisms elucidation may be performed by following the relaxation of the excited states over time with ultra-fast NIR spectroscopy. Alternatively, the effect dependence on concentration of Cu, and rare-earth ions in the perovskite can be investigated, which is a work in progress. The quenching of the chemically excited states may be further hindered by sample geometry and heat control. The experiments at hand gave indication of this. However, powder samples are suboptimal for this and need to be replaced by defined thin films.

Figure 3 Near infrared emission during catalytic combustion of methanol. The inset gives the flows of methanol-saturated air to pure air mixture changed in consecutive order (indicated by arrows). Note that the oxygen to methanol ratio, $r$, is decisive for the emission spectrum, while the absolute flows are less relevant.
Figure 4. Left panel: Simplified scheme of cataluminescence from methanol combustion on Er$_{0.15}$La$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_{3-\delta}$, in rich fuel-air mixtures (reducing conditions), and in rich fuel-air mixtures (oxidizing conditions), and in rich fuel-air mixtures (reducing conditions). The right figure compares the corresponding NIR and black-body spectra, alongside possible transitions in Er$^{3+}$ ions. ET denotes energy transfer from the catalytically active sites (Cu, golden spheres embedded in the perovskites) to the Er ions (blue). Plasmonic Cu nanoparticles (pyramid made of golden spheres) quench the luminescence.

4. Conclusions

Aim of this study was to alter the optical properties of a catalyst to suppress the black-body radiation and generate NIR in a selected energy range during catalytic methanol combustion. We demonstrated the selective IR-emission from Er-substituted SrTiO$_3$ during catalysis. Careful post-mortem as well as operando characterizations of quaternary perovskite Er$_{0.15}$La$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_{3-\delta}$ brought to fore the microstructural changes upon reaction conditions. The results provide arguments for cataluminescence, that is, non-thermal chemical energy conversion to light. The fraction of light with an energy above 0.8 eV generated by cataluminescence is markedly higher than obtained by black-body radiation making it attractive for future energy conversion devices. However, the selective NIR-emission lies in the range suboptimal for adjacent conversion to electricity in Si-based photovoltaic devices. Further work will thus be needed to utilize rare-earth ions with a better suited emission energy, and/or upconversion to the desired wavelengths. Research is needed to elucidate the role of Cu nanoparticles, and to optimize their impact.

5. Experimental Section

Optical Cell for Emission Measurements and Operando UV–Vis Diffusive Reflectance Spectroscopy: The aim of this study is to provide a quantitative estimate of the selective light emission from a catalyst during catalytic combustion (cataluminescence). For this a setup was developed, in which the catalyst was exposed to methanol air mixture, while the emitted NIR radiation was measured through a CaF$_2$ window. Figure 5 shows the employed optical setup consisting of a modified DRIFTS cell (Harrick cell HVC-DRP-3). The sample (powder) was placed over a Ni mesh in a heated metal cell to allow the flow of reactants and products through it. A methanol in air feed (0.5–2.0% vol.) was prepared by mixing two air streams, one of which was passed through a methanol saturator (23 °C corresponding to a methanol vapor pressure of 13 mbar). The two flow streams were independently regulated by manual flow controllers. Typical flows were 70 mL min$^{-1}$ methanol/air plus 30 mL min$^{-1}$ air for oxidizing conditions with a sample weight of 50 mg. The light emitted from the sample due to cataluminescence or thermoluminescence was collected by a fiber coupler (Thorlabs HCA 3) and coupled into a low OH fiber (Ocean Insight QP400-2-VIS-BX, transmitting best from 400–2100 nm) to an infrared (NIR) spectrometer (Ocean Optics FLAME-NIR-INTSMA25, grating N33: 150 lines per mm) for emission measurements. For UV–vis diffuse reflectance measurements, the sample was illuminated through the same fiber coupler using a split fiber optic cable (Ocean Optics, CUSTOM-BIF115765). The light source was a halogen lamp with fiber output, the spectrometer a UV–vis fiber spectrometer (Ocean Optics HR4000).

For calibration of the IR-emission signal, the cell was filled with a La$-$substituted SrTiO$_{3-\delta}$ with microstructural and optical properties similar to the Er-pervoskites. The cell was then heated to various temperatures at low air flows (10 mL min$^{-1}$) to reduce the temperature gradient in the cell. The measured spectrometer signal was then compared to the expected black body radiation at the corresponding temperatures. One proportionality factor (scaling factor) was sufficient to match measured spectra and black-body model (Figure 5). Using this factor, a universal calibration curve was derived correcting the small deviations from a non-ideal transmission function of the fiber and spectrometer. The thus derived luminescence (both thermal as well as cataluminescence) signal $I_{\text{lum}}$ was measured at various methanol to air mixtures.

Operando UV–vis diffusive reflectance spectroscopy were limited to the measurement of semiquantitative absorbance difference spectra $\Delta A_{\text{UV–vis}}$ with

$$\Delta A_{\text{UV–vis}} = -\ln \frac{I_{\text{UV–vis}}(f_o') - I_{\text{lum}}(f_o) - (I_{\text{UV–vis}}(f_i) - I_{\text{lum}}(f_i))}{I_{\text{UV–vis}}(f_o') - I_{\text{lum}}(f_o') + (I_{\text{UV–vis}}(f_i) - I_{\text{lum}}(f_i))} \quad (2)$$

The cataluminescence signal, $I_{\text{lum}}$, was subtracted as it also changes in the UV–vis spectral range when going from oxidizing to reducing flow mixtures $(f_o', f_i')$. 

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Figure 5. Left: Sketch of the optical cell used both for measuring the emission from the sample during catalysis and UV–vis diffuse difference spectra. For cataluminescence measurements, an NIR-spectrometer is attached to the fiber coupler (NIR light indicated by red arrows); for UV–vis diffuse reflectance spectroscopy, a split fiber is attached through which UV–vis light is brought to the sample, and the reflected part measured by a UV–vis spectrometer (light in and out indicated by yellow arrows). The right panel shows the calibration measurement making use of the black-body radiation from a similar but not luminescent sample La$_{0.3}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_3$.$\delta$.

Figure 6. Left panel: SEM micrograph (15 keV acceleration voltage, backscattered electron detector) of the calcinated powder showing the typical porous structure as result of the citrate preparation method. Right panel: XRD analysis of as prepared ErLaSTO (after calcination) and after hydrogen reduction at 200 and 850 °C. All peaks can be assigned to the four phases added from the ICSD reference. The oxide phases hardly change upon reduction, in contrast to the Cu phase, which is clearly visible only after reduction at higher temperatures (inset).

For product gas analysis, the exhaust of the Harrick cell was attached to infrared spectrometer with a 6 cm gas cell (Agilent 600 Series FTIR, Cary 600).

Post mortem UV–vis NIR diffuse reflectance spectra of the powders were measured with Jasco V770 spectrophotometer using powder holder (PSH-002, Jasco) in the back of the integrating sphere ISN-923 (Jasco, Easton, MD, USA). A reflectance measurement adapter was applied during the measurement to remove the specularly reflected light component. The reference of BaSO$_4$ powder used as an etalon white standard for a baseline correction and acetylene carbon black was used for a dark signal correction.

Materials Preparation and Characterization: La$_{0.15}$Er$_{0.15}$Sr$_{0.55}$Ti$_{0.95}$Cu$_{0.05}$O$_3$.$\delta$ (ErLaSTO) samples were prepared via an in-house adapted citrate-gel method$^{[12]}$: lanthanum(III) nitrate hexahydrate (99%, abcr GmbH), erbium(III) nitrate pentahydrate (99.9%, Acros organics), strontium(II) nitrate (99%, Acros organics), copper(II) nitrate trihydrate (98%, Carl Roth), and titanium diisopropoxide bis-acetylacetonate (Puriss, TBAD, Mateck, Germany) were used as starting precursors. Stoichiometric amounts of the metal salts were dissolved in an aqueous solution of citric acid (CA, 99.9%, VWR) and ethylene diamine tetra-acetic acid (EDTA, 99%, Carl Roth). An ammonia solution (25%, Merck) was used as a pH controlling agent. The ratio of metal ions to CA to EDTA was 1 : 1.5 : 1. After stirring, the precursor solution was placed in a drying furnace for 12 h at 140 °C. The resulting foam was crushed and heat treated for 6 h at 750 °C. Figure 6 is a micrograph showing the...
Figure 7. Left panel: Hydrogen signal indicative of the hydrogen consumption by the ErLaSTO during heating in Ar/H$_2$ mixture (TPR). The negative peaks mark the reduction process which is associated with the exsolution of Cu as shown by the increase of the Cu surface concentration (blue dots) derived from XPS measurements. Right panel: Selected XPS peaks (Cu 2p, Er 4d, and Ti 2p) of pristine, and reduced samples. The increase of the Cu surface concentration as deduced from the relative intensity ratios is reflected by a small binding energy shift of the Cu 2p peaks.

Cu nanoparticles were covered with an oxide skin resulting in a binding energy shift of 0.4 eV lower than expected from literature ($\Delta E_{\text{Cu-CuO}} = 1.4$ eV, $\Delta E_{\text{Cu-Cu$_2$O}} = 0.1$ eV, see Figure 7).

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Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are openly available at: http://doi.org/10.5281/zenodo.5055827.

Keywords

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[1] L. R. Kramer, A. L. O. Maran, S. S. De Souza, O. H. Ando, Energies 2019, 12, 16.
[2] G. Torsello, M. Lomascolo, A. Licciulli, D. Diso, S. Tundo, M. Mazzer, Nat. Mater. 2004, 3, 632.
[3] H. F. Ivey, J. Lumin. 1974, 8, 271.
[4] M. Breysse, B. Claudel, L. Faure, M. Guenin, R. J. Williams, T. Wolkenstein, J. Catal. 1976, 45, 137.
[5] L. Zhang, H. Song, Y. Su, Y. Lv, TrAC, Trends Anal. Chem. 2015, 67, 107.
[6] S. Wang, Z. Yuan, L. Zhang, Y. Lin, C. Lu, Analyst 2017, 142, 1415.
[7] J. Hu, L. Zhang, Y. Su, Y. Lv, Luminescence 2020, 35, 1174.
[8] S. Cimino, A. Gambirasi, L. Lisi, G. Mancino, M. Musiani, L. Vázquez-Gómez, E. Verlato, Chem. Eng. J. 2016, 285, 276.
[9] X. Wang, C. Zhang, Z. Zeng, Z. Hu, Combust. Sci. Technol. 2012, 184, 842.
[10] P. A. Deshpande, G. Madras, Phys. Chem. Chem. Phys. 2011, 13, 708.
[11] F. You, A. J. Bos, Q. Shi, S. Huang, P. Dorenbos, Phys. Rev. B 2012, 85, 115101.
[12] D. Burnat, R. Kontic, L. Holzer, P. Steiger, D. Ferri, A. Heel, J. Mater. Chem. A 2016, 4, 11939.
[13] W. Qi, C. Ruan, G. Wu, Y. Zhang, Y. Wang, K. Xie, Y. Wu, Int. J. Hydrogen Energy 2014, 39, 5485.
[14] J. Zhang, M.-R. Gao, J.-L. Luo, Chem. Mater. 2020, 32, 5424.
[15] A. A. Lyapin, S. V. Gushchin, S. V. Kuznetsov, P. A. Ryabochkina, A. S. Ermakov, V. V. Proydakova, V. V. Voronov, P. P. Fedorov, S. A. Artemov, A. D. Yapruntsev, V. K. Ivanov, Opt. Mater. Express 2018, 8, 1863.
[16] D. Avram, I. Tiseanu, B. Vasile, M. Florea, T. Carmen, Sci. Rep. 2018, 8, 18033.
[17] A. Pestryakov, V. Petranovskii, A. Kryazhov, O. Ozhereliev, N. Pfänder, A. Knop-Gericke, Chem. Phys. Lett. 2004, 385, 173.
[18] S. Nilsson, D. Albinsson, T. J. Antosiewicz, J. Fritzschke, C. Langhammer, Nanoscale 2019, 11, 20725.
[19] T. Ghodselahi, M. Vesaghi, Physica B 2011, 406, 2678.
[20] I. Yaremczuk, S. Meskinis, T. Bulavinets, A. Vasiliauskas, M. Andrulevicius, V. Fitio, Y. Bobitskia, S. Tamulevičius, Diamond Relat. Mater. 2019, 99, 107538.
[21] J. A. Jiménez, ChemPhysChem 2019, 20, 399.
[22] A. M. Stoneham, Rep. Prog. Phys. 1981, 44, 1251.
[23] C. Bettinali, G. Ferreresso, J. W. Manconi, J. Chem. Phys. 1969, 50, 3957.
[24] T. V. Perevalov, D. V. Gulyaev, V. S. Aliev, K. S. Zhuravlev, V. A. Gritsenko, A. P. Yelisseyev, J. Appl. Phys. 2014, 116, 244109.
[25] L.-J. Lai, H.-S. Sheu, Y.-K. Lin, Y.-C. Hsu, T.-C. Chu, J. Appl. Phys. 2006, 100, 103508.
[26] F. M. Pinto, V. Y. Suzuki, R. C. Silva, F. A. La Porta, Front. Mater. 2019, 6, 260.
[27] P. A. W. van der Heide, Surf. Interface Anal. 2002, 33, 414.
[28] D. Tahir, S. Tougaard, J. Phys.: Condens. Matter 2012, 24, 175002.