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Abstract. Latest results on $\beta\beta^0\nu$, $\beta\beta^0\nu\chi^0$ and $\beta\beta^2\nu$ decays of different isotopes from NEMO-3 double beta decay experiment are presented. In particular, new limits on neutrinoless double beta decay of $^{100}$Mo and $^{82}$Se have been obtained, $T_{1/2} > 4.6 \times 10^{23}$ y and $T_{1/2} > 1 \times 10^{23}$ y (90% C.L.), respectively. A possible next step with SuperNEMO detector is discussed.

1. Introduction
NEMO-3 detector is currently operating in the Frejus Underground Laboratory (4800 m w.e.). We present here recent results on $\beta\beta^0\nu$, $\beta\beta^0\nu\chi^0$ and $\beta\beta^2\nu$ transitions. At the same time the Collaboration has started to develop a SuperNEMO detector for 100 kg of $^{82}$Se.

2. The NEMO-3 experiment
2.1. The NEMO-3 detector
The experiment is based on the direct detection of two electrons produced by a double beta decay. NEMO-3 is able to operate with several double beta decay isotopes in the form of foils $\sim 50$ $\mu$m thick. The NEMO-3 detector accommodates $\sim 10$ kg of isotopes: 6.9 kg of $^{100}$Mo, 0.93 kg of $^{82}$Se, 0.4 kg of $^{116}$Cd, 0.45 kg of $^{130}$Te, 37 g of $^{150}$Nd, 9 g of $^{96}$Zr and 7 g of $^{48}$Ca. Particle detection in NEMO-3 consists of two parts. A tracking volume allows reconstruction of the tracks of charged particles inside the detector and a calorimeter measures the energy of $e^-$, $e^+$ and $\gamma$ particles. The tracking part is composed of 6180 drift cells, operating in Geiger mode, which provide three dimensional tracks. The tracking volume is surrounded by the calorimeter which is made of 1940 blocks of plastic scintillators. In addition, a magnetic field of 25 Gauss parallel to the detectors axis is created by a solenoid wound around the detector. The detector is surrounded by a passive shield.

The main characteristics of the detector’s performance are the following. The energy resolution of the scintillation counters lies in the interval of 14-17% (FWHM for 1 MeV electrons). The time resolution is 250 ps for an electron energy of 1 MeV. The reconstruction accuracy of a two electron (2e) vertex is around 1 cm. The characteristics of the detector are studied in special calibration runs with radioactive sources. A detailed description of the detector and its characteristics is presented in [1].
Figure 1. Spectra of the energy sum of the two electrons in the $\beta\beta$ energy window after 389 effective days of data collection from February 2003 until September 2004 (Phase I): (a) with 6.914 kg of $^{100}$Mo; (b) with 0.932 kg of $^{82}$Se; (c) with Copper and Tellurium foils. The shaded histograms are the expected backgrounds computed by Monte-Carlo simulations: dark (blue) is the $\beta\beta 2\nu$ contribution and light (green) is the Radon contribution. The solid line corresponds to the expected $\beta\beta 0\nu$ signal if $T_{1/2}(\beta\beta 0\nu) = 5 \times 10^{22}$ y.

2.2. Results

Measurement of the two neutrino double beta decay. For $^{100}$Mo and $^{82}$Se 389 effective days of data were analyzed. In case of $^{116}$Cd, $^{96}$Zr and $^{150}$Nd 168.4 days data were used. The results of the measurement are presented in Table 1.

| Nuclei   | Number of events | S/B ratio | $T_{1/2}, \text{y}$          |
|----------|------------------|-----------|-----------------------------|
| $^{100}$Mo | 219000           | 40        | $7.11 \pm 0.02(\text{stat}) \pm 0.54(\text{syst}) \times 10^{18}$ |
| $^{82}$Se  | 2750             | 4         | $9.6 \pm 0.3(\text{stat}) \pm 1.0(\text{syst}) \times 10^{19}$     |
| $^{116}$Cd | 1371             | 7.5       | $2.8 \pm 0.1(\text{stat}) \pm 0.3(\text{syst}) \times 10^{19}$     |
| $^{96}$Zr  | 72               | 0.9       | $2.0 \pm 0.3(\text{stat}) \pm 0.2(\text{syst}) \times 10^{19}$     |
| $^{150}$Nd | 449              | 2.8       | $9.7 \pm 0.7(\text{stat}) \pm 1.0(\text{syst}) \times 10^{18}$     |

Search for neutrinoless double beta decay. Figure 1(a) and (b) show the tail of the two-electron energy sum spectrum in the $\beta\beta 0\nu$ energy window for $^{100}$Mo and for $^{82}$Se respectively. The number of $2e^-$ events observed in the data is in agreement with the expected number of events from $\beta\beta 2\nu$ decay and simulations of the radon background.

In order to make the optimum use of all the information from the NEMO 3 detector, a maximum likelihood analysis has been applied to the $2e^-$ event sample above 2 MeV using the three available variables: the energy sum ($E_{\text{tot}}$) of the two electrons, the energy of each electron ($E_{\text{min}}$ is the minimum electron energy) and the angle between the two tracks ($\cos\theta$). With 389 effective days of data collection, limits at 90% C.L. obtained with the likelihood analysis are $T_{1/2}(\beta\beta 0\nu) > 4.6 \times 10^{23}$ years for $^{100}$Mo and $1.0 \times 10^{23}$ years for $^{82}$Se. The corresponding upper limits for the effective Majorana neutrino mass range from 0.7 to 2.8 eV for $^{100}$Mo and 1.7 to 4.9 eV for $^{82}$Se depending on the nuclear matrix element calculation [2, 3, 4, 5, 6]. For the hypothesis of a right-handed weak current, the limits at 90% C.L. are $T_{1/2}(\beta\beta 0\nu) > 1.7 \times 10^{25}$ years.
for $^{100}\text{Mo}$ and $0.7 \times 10^{23}$ years for $^{82}\text{Se}$, corresponding to an upper limit on the coupling constant of $\lambda < 2.5 \times 10^{-6}$ for $^{100}\text{Mo}$ and $3.8 \times 10^{-6}$ for $^{82}\text{Se}$ using the nuclear calculations from references [7, 8].

Decay with Majoron emission. In this case, the analysis of 8023 hours of NEMO-3 data is presented. The limits were obtained by analyzing the deviation in the shape of the energy distribution of the experimental data in comparison with calculated spectrum for $\beta\beta\nu$ decay. A maximum likelihood analysis was applied and different Majoron modes were investigated. The half-life limits for $^{100}\text{Mo}$ and $^{82}\text{Se}$ for the different decay modes are presented in Table 2. New limits on coupling constant of Majoron to neutrino were obtained. In particular, new limits on "ordinary" Majoron (spectral index 1) decay of $^{100}\text{Mo}$ and $^{82}\text{Se}$ correspond to bounds of $\langle g_{ee} \rangle < (0.4 - 1.8) \times 10^{-4}$ and $< (0.66 - 1.9) \times 10^{-4}$ using nuclear matrix element calculations from [2, 3, 4, 5].

Table 2. Limits at 90\% C.L. on $T_{1/2}(y)$ for different modes of double beta decay with Majoron emission. The "spectral index" $n$ defines the summed energy spectrum of the emitted electrons.

| Nuclei | $n = 1$ | $n = 2$ | $n = 3$ | $n = 7$ |
|--------|---------|---------|---------|---------|
| $^{100}\text{Mo}$ | $> 2.7 \times 10^{22}$ | $> 1.7 \times 10^{22}$ | $> 1.0 \times 10^{22}$ | $> 7 \times 10^{19}$ |
| $^{82}\text{Se}$ | $> 1.5 \times 10^{22}$ | $> 6.0 \times 10^{21}$ | $> 3.1 \times 10^{21}$ | $> 5.0 \times 10^{20}$ |

3. SuperNEMO experiment

The NEMO Collaboration is currently planning a future new, bigger detector. The main idea is to use the same experimental technique as in NEMO-2 [9] and NEMO-3 [1] experiments and to study 100 kg of $^{82}\text{Se}$. The planar geometry and modular scheme is proposed. The energy resolution will be improved up to $\sim 8 - 10\%$ (FWHM) at 1 MeV and the efficiency for $0\nu$ decay will be increased up to $\sim 20 - 40\%$. Other parameters will be the same as in case of NEMO-3. Sensitivity of the new experiment is estimated as $\sim 2 \times 10^{26} y$ for half-life or $\sim (0.04-0.1) eV$ for effective Majorana neutrino mass. A more detailed description of the SuperNEMO detector and its characteristics is presented in [10, 11].

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References

[1] Arnold R et al 2005 Nucl. Instrum. Methods A 536 79
[2] Rodin V A et al 2003 Phys. Rev. C 68 044302 (Preprint nucl-th/0503063)
[3] Simkovic F et al 1999 Phys. Rev. C 60 055502
[4] Civitarese O and Suhonen J 2003 Nucl. Phys. A 729 867
[5] Steica S and Klapdor-Kleingrothaus H V 2001 Nucl. Phys. A 694 269
[6] Caurier E et al 1996 Phys. Rev. Lett. 77 1954
[7] Aunola M and Suhonen J 1998 Nucl. Phys. A 643 207
[8] Suhonen J 2002 Nucl. Phys. A 700 649
[9] Arnold R et al 1995 Nucl. Instrum. Methods A 354 338
[10] Barabash A S 2002 Czech. J. Phys. 52 575
[11] Barabash A S 2004 Phys. At. Nucl. 67 1984