Stellar haloes of disc galaxies at $z \sim 1$

Ignacio Trujillo$^{1,2,*}$ and Judit Bakos$^{1,2}$

$^1$Instituto de Astrofísica de Canarias, Via Láctea s/n, 38200 La Laguna, Tenerife, Spain
$^2$Departamento de Astrofísica, Universidad de La Laguna, E-38205 La Laguna, Tenerife, Spain

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ABSTRACT

Taking advantage of the ultradeep near-infrared imaging obtained with the Hubble Space Telescope on the Hubble Ultra Deep Field, we detect and explore for the first time the properties of the stellar haloes of two Milky Way-like galaxies at $z \sim 1$. We find that the structural properties of those haloes (size and shape) are similar to the ones found in the local Universe. However, these high-$z$ stellar haloes are approximately three magnitudes brighter and exhibit bluer colours ($g-r \lesssim 0.3$ mag) than their local counterparts. The stellar populations of $z \sim 1$ stellar haloes are compatible with having ages $\lesssim 1$ Gyr. This implies that the stars in those haloes were formed basically at $1 < z < 2$. This result matches very well the theoretical predictions that locate most of the formation of the stellar haloes at those early epochs. A pure passive evolutionary scenario, where the stellar populations of our high-$z$ haloes simply fade to match the stellar halo properties found in the local universe, is consistent with our data.

Key words: galaxies: evolution – galaxies: formation – galaxies: haloes – galaxies: photometry – galaxies: spiral – galaxies: structure.

1 INTRODUCTION

A natural prediction of the hierarchical galaxy formation scenario is the ubiquitous presence of stellar haloes surrounding the galaxies (e.g. Eggen, Lynden-Bell & Sandage 1962; Searle & Zinn 1978; Steinmetz & Müller 1995; Bekki & Chiba 2001; Samland & Gerhard 2003). These extended and diffuse stellar components are formed from the debris of disrupted satellites accreted along the cosmic time (e.g. Brook, Kawata & Gibson 2004; Bullock & Johnston 2005; Abadi, Navarro & Steinmetz 2006). Recent cosmological simulations predict that the amount of stellar mass in these haloes should be $\sim 10^8-10^9 M_{\odot}$ for Milky Way (MW)-like objects. If these simulations are correct, most of the stellar halo mass would be assembled before $z \sim 1$, and we would expect since then a simple passive evolution of the stellar populations of these haloes towards the present (e.g. Cooper et al. 2010; Font et al. 2011).

From the observational point of view, the detection of stellar haloes is a major challenge. For example, Morrison (1993) estimates the surface brightness of the Galactic halo at the solar radius to be $V \sim 27.7$ mag arcsec$^{-2}$. In addition to the well-characterized MW halo (see a recent review by Helmi 2008), deep surveys of M31 (e.g. Ferguson et al. 2002; Irwin et al. 2005; Kalirai et al. 2006; Ibata et al. 2007; McConnell et al. 2009) have revealed an extensive halo (to $\sim 150$ kpc) with abundant substructure. M33 seems also to have a halo similar to the MW and M31, despite its smaller total mass (McConnachie et al. 2006). Beyond the Local Group, there is growing observational evidence showing that stellar haloes are ubiquitous and diverse. These studies have been conducted in nearby galaxies by using resolved star count techniques (e.g. Mouhcine et al. 2005; Mouhcine, Rejkuba & Ibata 2007; Mouhcine, Ibata & Rejkuba 2010; Ibata, Mouhcine & Rejkuba 2009; Radburn-Smith et al. 2011) or by extremely deep integrated photometry observations able to detect low surface brightness features (e.g. Sackett et al. 1994; Shang et al. 1998; Martínez-Delgado et al. 2008, 2009; Jablonka et al. 2010; Bakos & Trujillo 2012). Zibetti, White & Brinkmann (2004), see also Bergvall, Zackrisson & Caldwell (2010), stacking 1047 galaxies, were able to observe the properties of stellar haloes up to $z \sim 0.05$ ($\sim 200$ Mpc in distance). Finally, Zibetti & Ferguson (2004), using the Hubble Ultra Deep Field (HUDF; Beckwith et al. 2006), detected the stellar halo of a disc galaxy at $z = 0.32$. This last observation represents the current farthest detection of these faint components of the galaxies.

The observed number and physical properties of the stellar streams in nearby galaxies are in general agreement with the predictions from the cold dark matter model (e.g. Bell et al. 2008; Gilbert et al. 2009; McGaugh et al. 2009; Starkenburg et al. 2010). Even the large discrepancies in the properties of the stellar haloes of similar mass disc galaxies as the MW and M31 have been explained, within the cosmological context, as inherent to the system-to-system scatter in the halo formation histories of these objects (e.g. Font et al. 2011). To gain, consequently, a deeper understanding of the formation mechanisms of the disc galaxies is necessary: (a) to explore systematically the stellar haloes of nearby galaxies to even fainter surface brightness magnitudes ($V > 30$ mag arcsec$^{-2}$) where a plethora of substructures...
is expected to appear and (b) to probe the stellar haloes of high-redshift disc galaxies, where the properties of the stellar haloes are caught in earlier episodes of assembly. This paper deals with the second possibility. In particular, if the stellar halo formation scenarios are correct, a cosmic epoch that it is worth exploring is $z \sim 1$ where, as mentioned before, the haloes should be already in place.

If observing stellar haloes in the local universe is difficult, the observation of stellar haloes at $z \sim 1$ requires the deepest ever observations from the space. At $z = 1$, the cosmological dimming is $\sim 3$ mag arcsec$^{-2}$ and the optical red side of stellar emission (which contains most of the stellar flux) is only visible using near-infrared (NIR) data. Fortunately, the observations to start exploring the stellar haloes up to $z \sim 1$ have been recently taken by using the Wide Field Camera 3 (WFC3) onboard the Hubble Space Telescope (HST) imaging the HUDF. In this pilot paper, we present for the first time the stellar haloes of two MW-like galaxies at $z \sim 1$. We will show that these stellar haloes, in agreement with theoretical expectations, were much brighter and bluer than present-day stellar haloes of similar mass disc galaxies.

The structure of this paper is as follows. In Section 2, we describe the data used in this paper, the selection of our sample and the profiles extraction, and in Section 3 we present a detailed analysis of our results. Finally, in Section 4, we perform a short discussion and a summary of our work. The paper is complemented with an extensive Appendix where we expand on technical details used at conducting this paper.

All magnitudes in this paper are given in the AB system unless otherwise stated. Throughout we assume a flat-dominated cosmology ($\Omega_m = 0.3$, $\Omega_{\Lambda} = 0.7$ and $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$).

## 2 DATA AND SAMPLE SELECTION

Due to the faintness of the features we aim to explore in this paper, the study presented here requires the nowadays state-of-the-art imaging in terms of depth and resolution. This implies using the deepest optical and NIR imaging ever taken with the HST. The HUDF (RA = $03^h32^m39.0$, Dec. = $-27^\circ47^\prime29.1$ (J2000)) is the point on the sky with the best HST data for this project.

### 2.1 Data

In the optical regime, the HUDF has been observed (Beckwith et al. 2006) with the WFC of the Advanced Camera for Surveys (ACS) in four filters: F435W (B), F606W (V), F775W (i) and F850LP (z). The total exposure time in each band is as follows: 134 880 s (B), 135 320 s (V), 347 110 s (i) and 346 620 s (z), and the AB magnitude zero-points reached are: 25.673 (B), 26.486 (V), 25.654 (i) and 24.862 (z). Individual exposures in each filter were combined and drizzled to produce a single image per band with the following characteristics: 10 500 $\times$ 10 500 pixels with a pixel scale of 0.03 arcsec pixel$^{-1}$. These data were made public by the HUDF team at the following webpage: http://archive.stsci.edu/pub/hlsp/udf/acs-wfc/.

In the NIR, the HUDF was observed both by the Near Infrared Camera and Multi-Object Spectrometer (NICMOS) camera and more recently by the WFC3. The WFC3 observations are deeper and with higher resolution than NICMOS; however, the WFC3 observations do not cover entirely the HUDF field. For these reasons, in this paper, we have used both NIR data sets. The NICMOS HUDF (Thompson et al. 2005) data consist of F110W and F160W drizzled images of $3500 \times 3500$ pixels with a pixel scale of 0.09 arcsec pixel$^{-1}$. The zero-points are: 23.410 (F110W) and 23.220 (F160W). The images are available at http://archive.stsci.edu/pub/hlsp/udf/nicmos-treasury/.

The WFC3 HUDF imaging used here (Bouwens et al. 2011) only covers the very central (4.7 arcmin$^2$) part of the HUDF but to an extraordinary depth. The zero-points are: 26.27 (F105W), 26.25 (F125W) and 25.96 (F160W). The combined individual exposures in these filters were drizzled and have approximately $3000 \times 3000$ pixels with a pixel scale of 0.06 arcsec pixel$^{-1}$. The data are publicly available at: http://archive.stsci.edu/prepds/hudf09/.

### 2.2 Sample selection

The selection of our galaxies was conducted using the Rainbow Cosmological Database$^1$ published by Pérez-González et al. (2008); see also Barro et al. (2011a,b). This data base is a vast compilation of photometric and spectroscopic data for several of the deepest cosmological fields, such as The Great Observatories Origins Deep Survey (GOODS)-North and South, The Cosmological Evolution Survey (COSMOS) or the Extended Groth Strip, among others. Using all the available photometry, they have built spectral energy distributions (SEDs) covering the electromagnetic spectrum from the X-ray to the radio wavelengths. Analysing these SEDs, they have derived very robust photometric redshifts and accurate estimates of various stellar parameters (such as the stellar mass, the ultraviolet and infrared-based star formation rates, the stellar population age, etc.).

We have selected from the Rainbow data base those objects with spectroscopic redshift within 0.8 $< z < 1.2$ and $M_*$ $> 5 \times 10^8 M_\odot$ (Kroupa 2001 IMF). These constraints provide 15 galaxies. We have used the ACS $z$-band imaging to visually explore these objects. Two of these galaxies are ellipticals and are not considered in our study in what follows. From the remaining 13 galaxies, 9 objects are irregular or appear distorted by interactions. There are four bona fide spirals: UDF 3372, UDF 4438, UDF 5417 and UDF 9444. From those, we have finally selected only the two disc galaxies with the lowest inclinations. This mitigates potential effects of the dust on the surface brightness profiles we have obtained. In addition, while exploring the outer stellar haloes, less inclined galaxies are expected to be less affected by point spread function (PSF) effects (see de Jong 2008) than edge-on orientations. The characteristics of our galaxies are shown in Table 1. We have also shown these objects (in the $z$ and $H$ bands) in Fig. 1. Unfortunately, UDF 3372 is beyond the area covered by the HUDF WFC3 pointing, so we have had to rely our analysis on the NIR using the NICMOS data.

| Name    | RA (J2000) | Dec. (J2000) | Redshift | $M_{bul}$ (B band) ($\times 10^{10} M_\odot$) | Stellar mass | kpc arcsec$^{-1}$ |
|---------|------------|-------------|----------|---------------------------------|--------------|-----------------|
| UDF 3372 | 03$^h$32$^m$42$^s$3 | $-27^\circ$47$^\prime$46$^\prime$ | 0.996    | $-21.7$ | 2.4 | 8.001 |
| UDF 5417 | 03$^h$32$^m$39$^s$9 | $-27^\circ$47$^\prime$15$^\prime$ | 1.095    | $-21.6$ | 3.3 | 8.166 |

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$^1$ https://rainbowx.fis.ucm.es/
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Figure 1. Our sample of MW-like galaxies at $z \sim 1$: UDF 3372 and UDF 5417. Upper panels: the ACS $z$-band ($\sim B$-band rest frame) of the galaxies in our sample. Bottom panels: the NICMOS $F160W$ band (UDF 3372) and the WFC3 $F160W$ band (UDF 5417) ($\sim I$-band rest frame) of our galaxies. Listed in each figure is the galaxy name, its stellar mass and its spectroscopic redshift. The solid line indicates 1 arcsec angular size.

2.3 Profile extraction

We have extracted the surface brightness profiles in all the bands available for each of the galaxies. In the case of UDF 5417, we have obtained both WFC3 and NICMOS profiles, but we only show the deepest profiles corresponding to WFC3. The technique used to obtain the surface brightness profiles is fully explained in Bakos & Trujillo (2012). Summarizing, we extract radial surface brightness profiles on masked images in order to avoid contamination on our light profiles. We apply conservative masking on to sources which clearly do not belong to the galaxy, like foreground stars, background galaxies, etc. These sources are extracted by SExtractor (Bertin & Arnouts 1996). We use some SExtractor parameters, such as the measured flux, elongation, and similar, to determine the shape and size of these mask regions.

In order to extract radial surface brightness profiles representative of the most external part of the galaxy, we need to get characteristic values of the ellipticity and position angle of this region. We have done this by computing the second-order moment of the light distribution of the galaxy using the $z$-band image. The second-order moment is directly related to the position angle, the semimajor (A) and semiminor (B) axis lengths. We fix this ellipticity and position angle for all elliptical apertures. In each aperture, we estimate the galaxy flux by the $3\sigma$ rejected mean of the pixel values of that aperture. This helps to minimize the effect of morphological features like a spiral arm crossing the aperture.

Although our images are sky subtracted, our light profiles, however, can still be contaminated by some residual local sky background. This local background is estimated by using equally spaced apertures. We obtain the number count profiles of the galaxy up to very far distances, and we chose a large aperture where the profiles become flat, beyond the identifiable profiles of the galaxy. After a careful analysis of the profiles in all bands, the region where the sky is determined was chosen to be located at $55 < R < 75$ kpc (see Fig. A1). The $3\sigma$ rejected mean of the fluxes inside this aperture gives a robust estimate of the residual local sky background (for further information, see Pohlen & Trujillo 2006). The error of the background determination is key to decide the surface brightness level down to which we trust our profiles. Following a conservative approach, this is placed where the profiles obtained by either over- or undersubtracting the sky measurement by $\pm$σ start to deviate with more than 0.2 mag from the original profiles.

The observed profiles are shown in Fig. 2. Different vertical offsets have been applied to the profiles to facilitate the comparison among the different bands.

3 ANALYSIS

3.1 The shape of the surface brightness profiles

The two most remarkable features shown in the surface brightness profiles of Fig. 2 are the presence of a break in the disc at around 1 arcsec (i.e. $\sim 8$ kpc at $z = 1$), and some extra light above the outer disc expectations starting at $R \sim 2$ arcsec. These two features are observed in all the bands and in the two galaxies. This
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Figure 2. Observed profiles of the UDF 3372 (ACS and NICMOS) and UDF 5417 (ACS and WFC3) galaxies. The break features in the discs of the galaxies at ∼1 arcsec (i.e. ∼8 kpc) are observed in all the bands. Beyond $R = 2$ arcsec, in all the cases, there is an excess of light over the outer disc ($1 < R < 2$ arcsec) light exponential decrement. We identify this excess as the growing effect of the stellar halo in the surface brightness profiles of the galaxies.

strongly indicates that these features are real and are not connected to sky subtraction problems. We also note that they are observed in the NICMOS data (which are shallower than the WFC3), pointing out that (although the profiles are noisier) these data are still deep enough to detect the excess of light in the outer regions of the galaxies. We have checked that this is in fact the case using the NICMOS F160W filter of UDF 5417 and comparing this to the WFC3 F160W of the same object. The excess of light is also apparent in the NICMOS profile. The position of the break in the disc (at ∼1 arcsec) is pretty stable, independently of the different spatial resolution of the different bands used. We note, however, that in the very inner region of the UDF 3372, the presence of a bulge could be masked by the worst resolution of the NICMOS imaging compared to the WFC3 data.

We work in this paper under the hypothesis that the excess of light at $R > 2$ arcsec (i.e. $R > 16$ kpc) is due to the presence of stellar haloes in these objects. Our hypothesis is based on the similar shape of the surface brightness distribution of our galaxies compared to those found in nearby disc galaxies. In the local Universe, both based on integrated photometry (i.e. Bakos & Trujillo 2012) as well as using stellar counts (see e.g. Barker et al. 2012), the outer regions of low-inclined disc galaxies ($R > 10$ kpc) are characterized by an exponential surface brightness decrement followed by an excess of light over this exponential decay. This very outer excess of light is located at $R \gtrsim 20$ kpc and has surface brightness $\mu_R > 28$ mag arcsec$^{-2}$. These faint surface brightnesses are equivalent to the ones found in the literature using galaxies with edge-on orientations at similar radial distances above the galactic planes. For instance, the edge-on spiral NGC 4565 has, along its minor axis, a surface brightness of $\mu_{6660 \AA} = 27.5$ mag arcsec$^{-2}$ (i.e. $\mu_R \sim 28$ mag arcsec$^{-2}$ in the $R$ band) at 22 kpc above the disc (Wu et al. 2002). Similarly, Jablonka et al. (2010) also found that the stellar halo of the edge-on disc NGC 3957 has surface brightness of 28.5 mag arcsec$^{-2}$ in the $R$ band at 20 kpc over the disc plane. We conclude, consequently, that at distances of 20 kpc and surface brightness of $\mu_R > 28$ mag arcsec$^{-2}$ both in face-on and edge-on orientation of nearby disc galaxies, we are observing the same component of the galaxy, i.e. their stellar haloes. By analogy, we consider that the excess of light over the exponential disc decrement located in the peripheries (i.e. $R > 20$ kpc) of our high-$z$ face-on disc galaxies can be attributed to the stellar halo of these objects.

However, we should note, before claiming that the excess of light is our high-$z$ galaxies are only due to their stellar haloes, that de Jong (2008) showed that the effect of very extended PSF tails on the measurements of halo light can be significant (especially in the case of edge-on galaxies). de Jong conducted this study both in the Sloan Digital Sky Survey (SDSS) and the HUDF. Although de Jong (2008) concluded that in the face-on orientation, the effect of the PSF should not be so significant as in the edge-on case, it is worth exploring how much this effect can alter the estimation of...
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3.2 The effect of the PSF on the surface brightness profiles

To estimate the contribution from scattered light at large radii in our surface brightness profiles is necessary to model the intrinsic light distribution of the galaxies. We have done this by assuming that the galaxies are well described by a Sérsic \( r^{1/n} \) (1968) bulge plus a double-exponential disc. This accurately describes the inner (\( \lesssim 2 \) arcsec) profile. To calculate the observed distribution from the model distribution, the PSF needs to be determined out to a large radius (\( \sim 10 \) arcsec; de Jong 2008). Unfortunately, such an extended PSF cannot be accurately measured from the HUDF images itself as there are too few bright stars. Following de Jong (2008), we have therefore used the TINYTIM HST PSF modelling software to create artificial PSFs for each of our bands.

To illustrate the effect of the PSF, we show in Fig. 3 the model galaxy for UDF 5417 convolved with the PSF for the two reddest bands in both the ACS and in the WFC3 camera. For the rest of the bands and for galaxy UDF 3372, the PSF effect is shown in the Appendix. Clearly, a large fraction of the light seen in the surface brightness profile above the galaxy model is due to scattered light from the extended PSF. However, there is still a clear excess of light beyond 2 arcsec that cannot be explained by the PSF effect alone. It is necessary to add an extra component in the outer region of the galaxy to explain this excess of light.

In order to further investigate which is the origin of the outer light, we show in Fig. 4 the result of stacking the images from the three WFC3 bands of UDF 5417. This stacked image has been filtered with a median filter of 5 pixels width to enhance the lowest surface brightness features of the galaxy. The region corresponding to the excess of light above the outer disc expectations observed in the profiles (i.e. 16 < \( R < 40 \) kpc) are enclosed by two orange dashed circles. We can ascertain that this region is well beyond the main body (disc) of the galaxy. We also observe that the distribution of light in that region is not perfectly symmetrical. This absence of symmetry in a region beyond the disc could be an indication of ongoing accretion.

In the local Universe, deep surface brightness profiles (10 mag in range) of spiral galaxies, as the ones shown here for \( z = 1 \) objects, also exhibit the excess of light beyond the outer disc (see e.g. Bakos & Trujillo 2012). In present-day spirals, the surface brightness (\( \mu_0 \gtrsim 28 \) mag arcsec\(^{-2} \)), colours (\( g - r \gtrsim 0.6 \)) as well as the radial distances (\( R \gtrsim 15 \) kpc) of these light excesses correspond to the expected characteristics of stellar haloes.
We warn the reader about the dip present in the deconvolved (r band). In this band, the galaxies are brighter (\(z - r\) \(-1.09\)) galaxy. The image has been filtered with a median filter of 5 pixels width. These allow us to see the faintest features of the object. The dashed blue circle indicates the position of the break on the surface brightness profile of the disc of the galaxy. The region enclosed by the two orange dashed circles corresponds to the region of the break on the surface brightness profile of the disc of the galaxy. The region in this region are very uncertain.

3.3 Bulge–disc–halo decomposition

As we have seen in the previous section, the effect of the PSF is relevant in the outer regions of our galaxies and affects differently depending on the observed wavelength. To account for this effect, we have carried out a bulge–disc–stellar halo decomposition of our galaxies in all the observed bands. Those models allow us to work with the galaxies profiles without being affected by the PSF. From those PSF deconvolved models, we will obtain (see the next section) the rest-frame surface brightness profiles and the colour profiles of our objects not affected by the PSF. As we have done before, for illustrative purpose, we show the bulge–disc–stellar halo decomposition for the galaxy UDF 5417 in the F850LP and F160W WFC3 bands (Fig. 5) and leave the rest of the bands for the Appendix.

The structural parameters of the bulge–disc–stellar halo decomposition for each band are shown in Table 2. The structural decomposition is conducted on the 1D (i.e. 2D azimuthally averaged) profiles using our own code. Due to the large number of structural parameters involved in this decomposition, we have fixed several parameters. For instance, the size and shape of the bulge are based on the information retrieved from the F850LP band. In this band, the combination of depth and resolution for this structural component is the best that we have in our data. The magnitude of the bulge in the other wavelengths is then left free during the fitting process. Similarly, the shape (Gaussian) and size of the halo are also fixed and only the brightness is left free during the fit.

3.4 Rest-frame surface brightness and colour profiles

Using the observed surface brightness profiles, we have built the SEDs at different radial distances. We have linearly interpolated among the different observed filters to obtain the flux in the equivalent SDSS rest-frame bands: \(u\), \(g\) and \(r\). After doing this, we have corrected the rest-frame surface brightness profiles by adding the equivalent magnitude to the cosmological dimming at those redshifts. We show the rest-frame surface brightness profiles of our galaxies in Fig. 6. The rest-frame surface brightness profiles are used to build the rest-frame radial colour profiles as well.

The light in the outer regions of our galaxies is contaminated by the light scattered from the inner parts of the galaxy due to the extended PSF. The level of contamination is different depending on the wavelength. To account for that, both in the rest-frame surface brightness profiles as well as in the rest-frame colour profiles, we have used the bulge–disc–halo fit described in the above section. As we have done with the observed profiles, we have linearly interpolated among the different galaxy models of each observed filter to obtain the surface brightness of the galaxy model in the equivalent SDSS rest-frame band. The result of doing this is shown in Fig. 6. As we can see, the models reproduce fairly well both the rest-frame profiles and the rest-frame colours. Now, we repeat the exercise but this time using the PSF deconvolved models of each observed band (i.e. the model fits but without convolving them with the PSF). The new (’PSF effect free’) rest-frame surface brightness and colour profiles are shown in Fig. 6 with dashed lines.

At it is expected, the outer regions of the PSF deconvolved rest-frame surface brightness profiles are fainter than when the PSF effect is not accounted for. The effect on the \(u - g\) rest-frame colour is not very significant as this colour is mostly based on the observed ACS bands (which present the smaller PSFs). In the case of the \(g - r\) rest-frame colour, the change is larger since this colour is obtained from the information provided by the NIR cameras.

From Fig. 6, it follows immediately too that the surface brightness of our high-z galaxies are brighter (~2 mag arcsec\(^{-2}\)) than

Figure 4. Stacked image of all the WFC3 bands (F105W, F125W and F160W) of the UDF 5417 (\(z = 1.09\)) galaxy. The image has been filtered with a median filter of 5 pixels width. These allow us to see the faintest features of the object. The dashed blue circle indicates the position of the break on the surface brightness profile of the disc of the galaxy. The region enclosed by the two orange dashed circles corresponds to the region of the stellar halo explored in this paper. The visible structure of the halo is not perfectly symmetric suggesting some ongoing accretion.

2 Note that in the case of UDF 3372, it was not necessary to include a bulge to the decomposition as there was no evidence of any surface brightness increase in the centre of the profiles. To get our analysis as simple as possible, we just conducted a disc–stellar halo decomposition for this object.

We warn the reader about the dip present in the deconvolved \((g - r)\) colour between 10 to 20 kpc in the case of UDF 3372. We think this dip is artificial and not a real reflect of the underlying light distribution of the galaxy. This dip corresponds to the disc–stellar halo transition region and it is mainly produced by the shape of the deconvolved \(r\)-band profile for this galaxy. The rest-frame \(r\)-band mostly reflects the shape of the observed F110W band. Contrary to UDF 5417, where we are using much better resolution NIR data, in the case of UDF 3372 this radial region is strongly affected by the large wing of the NICMOS PSF (see e.g. Fig. B1) and consequently, the results in this region are very uncertain.
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**Figure 5.** Bulge–disc–stellar halo decomposition of the galaxy UDF 5417 in the reddest bands of the ACS ($F850LP$) and the WFC3 ($F160W$). The bulge, disc, halo as well as the galaxy models plus the galaxy models convolved with the PSFs are shown. The dashed lines in both panels indicate the limiting surface brightness down to which the surface brightness profiles are reliable.

| Band | $m_{\text{bulge}}$ | $r_e, \text{bulge}$ (arcsec) | $m_{\text{disc}}$ | $r_e, \text{disc}$ (arcsec) | $h_e, \text{break}$ (arcsec) | $h_e, \text{disc}$ (arcsec) | $m_{\text{halo}}$ | $r_e, \text{halo}$ (arcsec) |
|------|-------------------|-------------------------------|-----------------|-----------------------------|-----------------------------|-----------------------------|-----------------|-----------------------------|
| $b$  | 28.6              | 0.08                          | 1               | 23.10                       | 1.20                        | 0.90                        | 0.30            | 26.30                       | 2.5            |
| $v$  | 27.1              | 0.08                          | 1               | 22.68                       | 0.95                        | 0.90                        | 0.24            | 25.70                       | 2.5            |
| $i$  | 26.6              | 0.08                          | 1               | 22.01                       | 0.70                        | 0.90                        | 0.22            | 25.60                       | 2.5            |
| $z$  | 25.7              | 0.08                          | 1               | 21.45                       | 0.68                        | 0.94                        | 0.23            | 24.80                       | 1.7            |
| $F105W$ | 25.5          | 0.08                          | 1               | 21.30                       | 0.55                        | 0.99                        | 0.22            | 24.60                       | 1.7            |
| $F125W$ | 25.3          | 0.08                          | 1               | 21.04                       | 0.50                        | 1.05                        | 0.20            | 24.70                       | 1.7            |
| $F160W$ | 25.0          | 0.08                          | 1               | 20.85                       | 0.48                        | 1.10                        | 0.19            | 23.95                       | 1.7            |

Table 2. Bulge-disc-halo decomposition of the selected galaxies.

those shown in locally equivalent stellar mass galaxies. This is expected due to the large star formation activity and younger stellar population of those objects $\sim 8$ Gyr ago. Fortunately for us, not only the disc of the galaxies are brighter but also the stellar haloes. This is the reason why we can observe them even considering that the surface brightness of these galaxies are severely affected by the cosmological surface brightness dimming at $z \sim 1$ ($\sim 3$ mag arcsec$^{-2}$). The radial colour profiles of our galaxies (more clear in the data
Figure 6. Upper panels: the $u$, $g$ and $r$-band equivalent rest-frame surface brightness profiles obtained by the interpolation of the observed band profiles. The solid lines are the rest-frame surface brightness profiles obtained by the interpolation of the model fits to the observed band profiles (see Fig. 5; violet line). The dashed lines are the rest-frame surface brightness profiles obtained by the interpolation of the PSF deconvolved model fits to the observed bands (see Fig. 5; red line). The rest-frame surface brightness profiles have been corrected by the cosmological dimming. Bottom panels: the $g-r$ and $u-g$ rest-frame colour profiles are shown. The solid and dashed lines are the colours obtained from the model fits to the data (solid line) and the PSF deconvolved models (dashed lines).

Table 3. Structural properties of the selected galaxies in the rest-frame SDSS $r$ band.

| Name       | $R_{\text{break}}$ (arcsec) | $R_{\text{break}}$ (kpc) | $\mu_r, \text{break}$ (mag arcsec$^{-2}$) | $R_e, \text{halo}$ (kpc) | $\langle \mu_e, \text{halo} \rangle$ (mag arcsec$^{-2}$) | $L_{\text{halo}}/L_{\text{total}}$ (r band) | $(g - r)_{\text{halo}}$ (mag) |
|------------|----------------------------|--------------------------|------------------------------------------|--------------------------|------------------------------------------------|------------------------------------------|----------------------------------|
| UDF 3372   | 0.95 ± 0.02                | 7.6 ± 0.1                | 20.0 ± 0.03                              | 19.6 ± 2.3               | 25.4 ± 0.25                                         | 0.040 ± 0.010                          | 0.05 ± 0.30                     |
| UDF 5417   | 1.05 ± 0.02                | 8.6 ± 0.1                | 20.3 ± 0.03                              | 13.6 ± 0.6               | 24.4 ± 0.10                                         | 0.041 ± 0.004                          | 0.15 ± 0.15                     |

itself than in the models) show the U-shape (associated with the break in the stellar disc) which has been observed in the nearby Universe (Bakos, Trujillo & Pohlen 2008) as well as in $z \lesssim 1$ disc galaxies (Azzollini, Trujillo & Beckman 2008a). In agreement with the fact that high-$z$ galaxies have younger stellar populations at $z = 1$ than now, the observed colours of our disc galaxies are bluer than in local objects. For instance, whereas in the local Universe, the surface brightness profiles usually have a $(g - r) = 0.47 ± 0.02$ mag at the break radius, our galaxies present $(g - r) < 0.3$ mag at this feature.

In Table 3, we provide some structural parameters of the disc component of the galaxies. In particular, we present the radial position of the disc break and the $r$-band rest-frame surface brightness (after correcting by the cosmological dimming) at this position. Both the values that we get for the surface brightness of the breaks ($\mu_r, \text{break} \sim 20.2$ mag arcsec$^{-2}$) and their radial positions ($R_{\text{break}} \sim 8.1$ kpc) are typical of the population of disc galaxies with similar absolute magnitudes and stellar mass at those redshifts (e.g. Azzollini, Trujillo & Beckman 2008b).

3.5 Stellar halo properties

Following a similar analysis to the one presented in Bakos & Trujillo (2012) for local galaxies and as we have done earlier, we characterize the structural properties of the stellar haloes at high-$z$ by decomposing the PSF deconvolved $r$-band surface brightness rest-frame profile into three different components: a Sérsic $r^{1/n}$ (1968) bulge, a double-exponential disc and a stellar halo described by a Sérsic model with $n = 0.5$ (see Fig. 7). This later choice seems somehow arbitrary but it has been chosen because it represents well the shape
of the profile in the outer region. It also allows a direct comparison with the analysis of the stellar haloes conducted at $z = 0$.

In Table 3, we provide the effective radii of the haloes and their mean surface brightness according to the rest-frame $r$ band. In the same table, we also show the fraction of light that is contained in the stellar haloes compared to the total light of the galaxies ($L_{\text{halo}}/L_{\text{total}}$). The effective radii of our haloes ($R_{\text{e, halo}} \sim 17$ kpc) are within the range of the ones found in the local Universe. The most important difference of the high-$z$ stellar haloes compared to the local ones is in connection with their mean effective surface brightness. At high-$z$, the values found ($\langle \mu_{e, \text{halo}} \rangle \sim 24.9$ mag arcsec$^{-2}$ ($r$ band) are significantly much brighter ($\gtrsim 3$ mag arcsec$^{-2}$) than those measured in nearby galaxies. We finally have obtained the $(g - r)$ colour at the stellar halo position by using the PSF deconvolved model colour profiles at $20 < R < 25$ kpc (see Fig. 6). Ideally, we would like to explore these colours farther away but the errors are so large that this estimation becomes prohibitive. For this reason, we concentrate on the above region. Nonetheless, if we trust the bulge–disc–halo decomposition conducted here, these colours should not be significantly affected by the colours of the outer disc population after the PSF effect correction. To be conservative, the halo colour error as well as the error on the mean effective surface brightness of the halo is based on the observed error measured in the surface brightness and colour distributions at the farthest halo regions we can explore with confidence (i.e. $20 < R < 25$ kpc). We acknowledge that these errors do not account for potential variations of the shape of the halo (which we have fixed here to a Gaussian $n = 0.5$ shape). So, we warn the reader that there is a potential source of systematic bias that cannot be easily quantified without having deeper and/or higher resolution images.

3.6 Comparison with the local sample

To put our high-$z$ galaxies in context, we have compared the properties of their stellar haloes with the ones found in the local Universe. The result of this comparison is presented in Fig. 8. In the left-hand panel, we have the light fraction of the stellar haloes (in the rest-frame SDSS $r$ band) versus the total galaxy stellar mass. We have taken the local values from different authors. These works are listed in the caption to Fig. 8. We find that the amount of light contained in the high-$z$ stellar haloes is very similar to the values found in our local reference sample ($\sim 4$ per cent). This result implies that, whatever the complex mechanisms acting on the evolution of the galaxies as a whole and the stellar haloes as a separate component, they should produce a decrease of the brightness in both structures which is similar in order to preserve such fraction constant with time.

In the right-hand panel of Fig. 8, we compare the mean effective surface brightness and the $(g - r)$ colour of the stellar haloes both at high and low-$z$. As we have been claiming through the paper, the stellar haloes at high-$z$ are significantly much brighter than today. Also, their colours are bluer ($\langle g - r \rangle \sim 0.1$ mag) than in the local universe ($\langle g - r \rangle \sim 0.7$ mag; see also Zibetti et al. 2004; Zackrisson et al. 2006; Bergvall et al. 2010). Can the high-$z$ stellar haloes in the high-$z$ Universe evolve into the ones that we observe locally by passive evolution? The most simple exercise that we can
conduct this question is to follow the tracks of single stellar populations (SSPs) with cosmic time and to probe whether the passive evolution of the stellar population in the haloes can reproduce the local properties. These tracks have been constructed using the Vazdekis et al. (2010) models for SSPs with two different IMFs (Salpeter 1955 and Kroupa 2001). We find that the passive evolution moves the colour and the surface brightness of the high-\(z\) stellar haloes towards the observed values in the local Universe. We explore the consequences of this finding in the next section.

4 DISCUSSION AND CONCLUSIONS

In the current paradigm, stellar haloes of galaxies have a dual nature. The outer halo formed primarily by accretion/mergers, whereas the inner halo is formed mainly through a dissipative collapse (as a result of gas-rich mergers at early times). Recent cosmological simulations (e.g. Cooper et al. 2010; Font et al. 2011; Tissera, White & Scannapieco 2012) indicate that both phenomena take place at \(z > 1\). If this is the case, one would expect that the stellar haloes will evolve since that epoch by mostly fading without any significant change in their structures. Our present observations allow us to check this important prediction of the disc galaxy formation scenario.

According to our data, the structural properties of our high-\(z\) stellar haloes are compatible with being at place at \(z \sim 1\). We base this assertion on the following: if we assume that the high-\(z\) haloes have the same shape than the local ones (i.e. if they can be well described with \(n = 0.5\) profiles as it seems to be the case), then the sizes of these components (\(R_t, halo \sim 17\) kpc) as well as their contribution to the total light of the galaxies (\(\sim 4\) per cent) are very similar to the ones found in the local Universe. The major difference with respect to the local galaxies is that our haloes are approximately three magnitudes brighter than their local counterparts. They also present bluer colours than the local stellar haloes. Both results agree with the idea that high-\(z\) stellar haloes are significantly much younger than the present-day ones. According to the SSP evolutionary tracks that we have explored, the colours of our high-\(z\) stellar haloes correspond to stellar populations with an age \(\lesssim 1\) Gyr. This implies that the stars in those haloes were formed basically at \(1 < z < 2\). This matches very well the theoretical predictions that move most of the formation of the stellar haloes at those epochs.

Our sample is so far very small (only two objects), but this pilot work can be easily extended towards other spiral galaxies in the HUDF at lower redshifts and with different inclinations. Our two disc galaxies have not been pre-selected to have bright (i.e. detectable) stellar haloes but we admit that we cannot infer, with only two objects, whether this is the general behaviour of the whole disc galaxy population at \(z \sim 1\). It will be interesting to explore in the future whether other disc galaxies have prominent stellar haloes as the ones discussed in this paper. In addition, it will be worth probing whether the properties of other high-\(z\) stellar haloes also fit within a scenario where a pure passive evolution of their stellar populations matures them towards the present-day Universe stellar haloes.

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APPENDIX A: MASKING

The characterization of the stellar haloes (e.g. shape, structural parameters and colours) can be contaminated by the presence of nearby sources like foreground stars, background galaxies, etc. In this Appendix, we show the masks that we have used in this paper to avoid this problem (see Fig. A1). We created two different mastermasks for each galaxy: one in the optical regime and another one in the NIR. The background contaminating sources are extracted by SExtractor from a master image composed of all the bands. In the case of the optical regime, this corresponds to the four ACS filters, and in the case of the NIR data, this master image is composed of the two NICMOS filters (for the galaxy UDF 3372) and of the three WF3C filters (for the galaxy UDF 5417). The need to separate between the optical and NIR regimes at creating the masks is justified by the significantly different PSFs sizes and cameras characteristics. Nonetheless, as we will see later, most of the regions masked in both regimes are in common being the NIR-masked areas larger than in the optical as expected due to their larger PSFs.

The master images are a stack of the different filters scaled to the z-band flux (optical) and H-band flux (NIR). Sources are extracted in the ‘cold mode’, optimized for brighter objects. We use some SExtractor parameters such as the measured flux, elongation and similar, to determine the shape and size of these mask regions. In some cases, problematic sources, saturated stars, for instance, are masked manually. Since the masking is based on information coming from all filters in each regime, we use the resulting masks as mastermasks, and it is applied to all filters in the optical regime (mastermask in the optical) and in the NIR (mastermasks in the NIR).

APPENDIX B: PSF EFFECT

We show in this section the effect of the PSF on the surface brightness profiles of both galaxies and in all the bands used in this work. For the galaxy UDF 3372, we show the effect of the PSF in Fig. B1 and for the galaxy UDF 5417 in Fig. B2.

It is worth noting that the excess of light above the PSF model prediction is observed in all the profiles except in the F160W NICMOS profile of the galaxy UDF 3372. The effect of the PSF in that band is so large that it does not permit exploring any potential stellar halo light at those distances. Also, in the F435W band, in both galaxies, the data are not deep enough to explore the presence of a stellar halo at those (ultraviolet rest-frame) wavelengths.
Figure A1. Regions masked (shown in black) in this work. For each regime (optical and NIR), we have created different masks to account for the significantly different PSFs and cameras characteristics. The region chosen to determine the sky is enclosed by two orange dashed circles.

Figure B1. Surface brightness profile of the UDF 3372 galaxy in all the observed bands. The PSF profiles normalized to the same central surface brightness of the galaxy, as well as the galaxy models plus the galaxy models convolved with the PSFs are shown. The dashed lines in both panels indicate the limiting surface brightness down to which the surface brightness profiles are reliable.
APPENDIX C: BULGE–DISC– STELLAR HALO DECOMPOSITIONS

In this Appendix, we show the bulge–disc–stellar halo decomposition for the two galaxies explored in this work in all the bands that we have used in this paper. The results of doing this decomposition are presented in Fig. C1 for UDF 3372 and Fig. C2 for UDF 5417.

APPENDIX D: RESULTS BASED ON MODELLING THE HALO WITH AN EXPONENTIAL PROFILE

In this work, we have assumed that the surface brightness distribution of the stellar halo follows a Gaussian distribution (i.e. a Sérsic model with an $n$ index fixed to $n = 0.5$). This is done for two reasons: first, it produces a fit in all the bands which agrees very well with the data distribution, and secondly, it allows us to compare with the results obtained in the local sample by Bakos & Trujillo (2012), where this model was used to extract the properties of the stellar haloes. Nonetheless, the exact determination of the shape of the stellar halo is uncertain due to the limited spatial range that we can cover with our observations (not reliable beyond 40 kpc). For that reason, we have repeated our full analysis based on an exponential distribution for the outer component. Some of the values obtained directly from this analysis cannot be straightforward compared with the local sample by Bakos & Trujillo (2012) as those values (for instance: the mean surface brightness, the effective radius of the halo, etc.) strongly depend on the model used to fit the outer component (see e.g. table 3 of Gilbert et al. 2012). However, other parameters such as the fraction of light in the stellar halo compared to the total luminosity of the galaxy and/or the colour-corrected profiles help to probe how dependent are our results on the chosen shape for the stellar outer component.

Exponential profiles for fitting the outer component of the stellar halo produce larger effective radii and fainter surface brightnesses [UDF 3372: $R_{e, \text{halo}} = 47.2 \pm 5.7$ kpc, $\langle \mu_e, \text{halo} \rangle = 26.4 \pm 0.3$ mag arcsec$^{-2}$ (r-band rest frame); UDF 5417: $R_{e, \text{halo}} = 44.1 \pm 2.2$ kpc, $\langle \mu_e, \text{halo} \rangle = 26.4 \pm 0.1$ mag arcsec$^{-2}$ (r-band rest frame)] than that in the case of Gaussian models. The stellar halo light contributions using exponential model rise to 9.1 $\pm$ 2.3 per cent (UDF 3372) and 6.5 $\pm$ 0.7 per cent (UDF 5417). This is understandable due to the larger amount of light that is located on the tail of an exponential model compared to a Gaussian distribution. Nonetheless, in the case of UDF 5417, where our data are most reliable and deep, the stellar halo light contribution using both models are relatively similar.

Finally, we have explored how the colour correction we have applied due to the effect of the PSF in our surface brightness profiles altered when we use an exponential model. In Fig. D1, we show the $u$, $g$ and $r$ rest-frame surface brightness profiles as well as the $u - g$ and $g - r$ colour PSF corrected radial profiles when using an exponential model to fit the outer region. This figure shows that, within the error bars, the PSF colour-corrected radial profiles are not significantly different assuming a Gaussian or an exponential model.
Figure C1. Bulge-disc–halo decomposition of the UDF 3372 galaxy in all the observed bands. The bulge, disc, halo as well as the galaxy models plus the galaxy models convolved with the PSFs are shown. The dashed lines in both panels indicate the limiting surface brightness down to which the surface brightness profiles are reliable.

Figure C2. Bulge-disc–halo decomposition of the UDF 5417 galaxy in all the observed bands. The bulge, disc, halo as well as the galaxy models plus the galaxy models convolved with the PSFs are shown. The dashed lines in both panels indicate the limiting surface brightness down to which the surface brightness profiles are reliable.
Figure D1. Upper panels: the $u$, $g$ and $r$-band equivalent rest-frame surface brightness profiles obtained by the interpolation of the observed band profiles. The solid lines represent the rest-frame surface brightness profiles obtained by the interpolation of the model fits to the observed band profiles. The dashed lines represent the rest-frame surface brightness profiles obtained by the interpolation of the PSF deconvolved model fits to the observed bands. The rest-frame surface brightness profiles have been corrected by the cosmological dimming. Bottom panels: the $g - r$ and $u - g$ rest-frame colour profiles are shown. The solid and dashed lines are the colours obtained from the model fits to the data (solid line) and the PSF deconvolved models (dashed lines). The stellar haloes have been modelled using exponential profiles.

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