Magnetic Geodesic Flows on Coadjoint Orbits

Alexey V. Bolsinov , Božidar Jovanović

Department of Mechanics and Mathematics, Moscow State University
119992, Moscow, Russia, e-mail: bolsinov@mech.math.msu.su

and

Mathematical Institute SANU
Kneza Mihaila 35, 11000 Belgrade, Serbia, e-mail: bozaj@mi.sanu.ac.yu

October 1, 2018

Abstract

We describe a class of completely integrable $G$-invariant magnetic geodesic flows on (co)adjoint orbits of a compact connected Lie group $G$ with magnetic field given by the Kirillov-Konstant 2-form.

1 Introduction

Let $Q$ be a smooth manifold with a local coordinate system $x^1,\ldots,x^n$ and Riemannian metric $g = (g_{ij})$. The inertial motion of the unit mass particle under the influence of the additional magnetic field given by a closed 2-form

$$\Omega = \sum_{1\leq i<j\leq n} F_{ij}(x)dx^i \wedge dx^j,$$

is described by the following equations on the phase space $T^*Q$:

$$\frac{dx^i}{dt} = \frac{\partial H}{\partial p_i}, \quad \frac{dp_i}{dt} = -\frac{\partial H}{\partial x^i} + \sum_{j=1}^n F_{ij} \frac{\partial H}{\partial p_j},$$

(1)

where $p_j = g_{ij}x^i$ are canonical momenta and the Hamiltonian is $H(x,p) = \frac{1}{2} \sum g^{ij}p_ip_j$. Here $g^{ij}$ are the coefficients of the tensor inverse to the metric.

The equations (1) are Hamiltonian with respect to the symplectic form $\omega + \rho^*\Omega$, where $\omega = \sum dp_i \wedge dx^i$ is the canonical symplectic form on $T^*Q$ and $\rho^* \Omega$ is the pullback of the magnetic 2-form by the evolution generated by the Hamiltonian $H$. The equations (1) can be written in the Hamiltonian form with respect to the symplectic form $\omega + \rho^*\Omega$.
\( \rho : T^*Q \to Q \) is the natural projection. Namely, the corresponding Poisson bracket is given by

\[
\{ f, g \} = \sum_{i=1}^n \left( \frac{\partial f}{\partial x^i} \frac{\partial g}{\partial p_i} - \frac{\partial g}{\partial x^i} \frac{\partial f}{\partial p_i} \right) + \sum_{i,j=1}^n F_{ij} \frac{\partial f}{\partial p_i} \frac{\partial g}{\partial p_j},
\]

and the Hamiltonian equations \( \dot{f} = \{ f, H \} \) read 11. The flow 11 is called magnetic geodesic flow on the Riemannian manifold \((Q, g)\) with respect to the magnetic field \(\Omega\).

In this paper we consider \(G\)-invariant magnetic geodesic flows on (co)adjoint orbits \(O(a)\) of a compact connected Lie group \(G\), where \(\Omega\) is the Kirillov-Konstant 2-form (Theorem 1). The non-commutative integrability of the systems, for the normal metrics, is proved recently by Efimov 5, 6. Following 3, we give a new, short proof of the non-commutative integrability (Theorem 2). In addition, the usual Liouville integrability by means of commuting analytic integrals is shown. One can use the commuting integrals in order to deform the normal metric to a certain class of \(G\)-invariant metrics on \(O\) with completely integrable magnetic geodesic flows as well (Theorem 3).

2 Magnetic Coadjoint Orbits

Let \(G\) be a compact connected Lie group with the Lie algebra \(\mathfrak{g} = T_e G\). Let us fix some \(\text{Ad}_G\)-invariant scalar product \(\langle \cdot, \cdot \rangle\) on \(\mathfrak{g}\). By the use of \(\langle \cdot, \cdot \rangle\) we identify \(\mathfrak{g}^*\) with \(\mathfrak{g}\).

Consider the adjoint action of \(G\) and the \(G\)-orbit \(O(a) = \text{Ad}_G(a)\) through an element \(a \in \mathfrak{g}\). Let \(\xi \in \mathfrak{g}\) and \(x = \text{Ad}_g(a)\). Since

\[
\xi_x = \frac{d}{ds} \text{Ad}_{\exp(s\xi)}(x)|_{s=0} = [\xi, x],
\]

the tangent space \(T_x O(a)\) is simply \([\mathfrak{g}, x]\), i.e., it is the orthogonal complement to \(\text{ann}(x) = \{ \eta \in \mathfrak{g} \mid [\eta, x] = 0 \}\). By definition, the Kirillov-Konstant symplectic form \(\Omega\) on \(O(a)\) is a \(G\)-invariant form, given by

\[
\Omega(\eta_1, \eta_2)|_x = -\langle x, [\xi_1, \xi_2]\rangle, \quad \eta_i = [\xi_i, x], \quad i = 1, 2.
\]

Similarly, the scalar product \(\langle \cdot, \cdot \rangle\) induces the normal metric \(K_0\) on \(O(a)\) as follows

\[
K_0(\eta_1, \eta_2)|_x = \langle \xi_1, \xi_2\rangle, \quad \eta_i = [\xi_i, x], \quad i = 1, 2.
\]

The cotangent bundle \(T^* O(a)\) can be realised as a submanifold of \(\mathfrak{g} \times \mathfrak{g}\)

\[
T^* O(a) = \{(x, p) \mid x = \text{Ad}_g(a), p \in \text{ann}(x)^\perp\},
\]

with the paring between \(p \in T_x^* O(a)\) and \(\eta \in T_x O(a)\) given by \(p(\eta) = \langle p, \eta \rangle\). Then the canonical symplectic form \(\omega\) on \(T^* O(a)\) can be seen as a restriction of
the canonical linear symplectic form of the ambient space $g \times g$: $\sum_{i=1}^{\dim g} dp_i \wedge dx_i$, where $p_i, x_i$ are coordinates of $p$ and $x$ with respect to some base of $g$.

The $G$-action

$$g \cdot (x, p) = (\operatorname{Ad}_g x, \operatorname{Ad}_g p)$$

is Hamiltonian on $(T^*O(a), \omega)$. From (3) we find that the momentum mapping is given by the relation $\langle \Phi_0(x, p), \xi \rangle = \langle p, \xi_x \rangle = \langle p, [\xi, x] \rangle$. That is $\Phi_0(x, p) = [x, p]$.

Following Efimov [5, 6], we consider magnetic geodesic flows on $O(a)$ with respect to the magnetic fields $\epsilon \Omega$, where $\Omega$ is Kirillov-Konstant 2-form [4] and $\epsilon \in \mathbb{R}$. According to (1), the adding of magnetic field $\epsilon \Omega$ to the system reflects as a perturbation of the system in $p$-variable by the magnetic force $\Pi_\epsilon$, determined by $\langle \Pi_\epsilon, \eta \rangle = -\epsilon \langle x, [\operatorname{ad}_x^{-1} \eta, \operatorname{ad}_x^{-1} \dot{x}] \rangle$, $\eta \in T_xO(a)$. Hence $\Pi_\epsilon = -\epsilon \operatorname{ad}_x^{-1} \dot{x}$.

The $G$-action (6) is Hamiltonian on $(T^*O(a), \omega + \epsilon \Omega)$ as well [6, 7]. In our notation we have that the momentum mapping reads

$$\Phi_\epsilon(x, p) = \Phi_0(x, p) + \epsilon x = [x, p] + \epsilon x.$$

**G-Invariant Magnetic Geodesic Flows.** The $G$-invariant metrics on $O(a)$ are in one-to-one correspondence with $\operatorname{Ad}_{Ga}$-invariant positive definite operators $\varphi: v \to v, \quad \operatorname{Ad}_g \circ \varphi = \varphi \circ \operatorname{Ad}_g, \quad g \in G_a,$

where $v = T_0O(a) = \operatorname{ann}(a)^\perp$ and $G_a$ is the isotropy group of $a$. Namely, for a given $\varphi$, we define

$$\varphi_x = \operatorname{Ad}_g \circ \varphi \circ \operatorname{Ad}_{g^{-1}}: T_xO(a) \to T_xO(a), \quad x = \operatorname{Ad}_g(a),$$

and a $G$-invariant metric $K_\varphi(\eta_1, \eta_2)|_x = \langle \varphi_x \eta_1, \eta_2 \rangle$. After Legendre transformation $T^*O(a) \to T^*O(a)$ with respect to $K_\varphi$, we get the Hamiltonian function for the given metric:

$$H_\varphi(x, p) = \frac{1}{2}\langle \varphi_x^{-1} p, p \rangle.$$

**Theorem 1** The equations of the magnetic geodesic flow on $(O(a), K_\varphi)$ with respect to the magnetic term $\epsilon \Omega$, in redundant variables $(x, p)$, are given by

$$\dot{x} = \varphi_x^{-1} p,$$

$$\dot{p} = \operatorname{ad}_x^{-1}[p, \varphi_x^{-1} p] - \operatorname{pr}_{\operatorname{ann}(x)}[\operatorname{ad}_x^{-1} \varphi_x^{-1} p, p] - \epsilon \operatorname{ad}_x^{-1} \varphi_x^{-1} p.$$  

In particular, the magnetic flow of the normal metric [3] reads

$$\dot{x} = [x, p],$$

$$\dot{p} = [x, p] + \epsilon [x, p]$$.
Proof. The equation (11) is just the inverse of the Legendre transformation. We can derive (8) simply by using the conservation of the momentum mapping $Φ_ε$ for $G$-invariant Hamiltonians. We have

\[
\frac{d}{dt} Φ_ε(x,p) = [\dot{x}, p] + [x, \dot{p}] + ε\dot{x} = 0 \\
= [φ_x^{-1}p, p] + [x, \dot{p}] + ε[x, ad_x^{-1}φ_x^{-1}p] = 0.
\]

(11)

Since $φ^{-1}$ is $Ad_{G_x}$-invariant, the term $[φ_x^{-1}p, p]$ belongs to $ann(x)^\perp$. Thus from (11) we get

$$pr_{ann(x)\perp} \dot{p} = ad_x^{-1}[p, φ_x^{-1}p] - εad_x^{-1}φ_x^{-1}p. \tag{12}$$

In order to find $pr_{ann(x)\perp} \dot{p}$, take the (local) orthonormal base $e_1(x), . . . , e_r(x)$ of $ann(x)$. Then $pr_{ann(x)\perp} \dot{p}$ is determined from the condition that the trajectory $(x(t), p(t))$ satisfies constraints

$$\frac{d}{dt} \langle p, e_i(x) \rangle = \langle \dot{p}, e_i(x) \rangle + \langle p, \dot{e}_i(x) \rangle = 0, \quad i = 1, . . . , r. \tag{13}$$

From $[e_i(x), x] ≡ 0, i = 1, . . . , r,$ we get

$$[\dot{e}_i(x), x] + [e_i(x), \dot{x}] = [\dot{e}_i(x), x] + [e_i(x), [x, ad_x^{-1}φ_x^{-1}p]] = 0 \quad i = 1, . . . , r. \tag{14}$$

Furthermore, combining (14) and the Jacobi identities

$$[e_i, [x, ad_x^{-1}φ_x^{-1}p]] + [x, [ad_x^{-1}φ_x^{-1}p, e_i]] + [ad_x^{-1}φ_x^{-1}p, e_i, x] = 0, \quad i = 1, . . . , r$$

we obtain $\dot{e}_i(x) = [e_i(x), ad_x^{-1}φ_x^{-1}p]$ (modulo $ann(x)$). Whence, using (13) we get $\langle \dot{p}, e_i(x) \rangle + \langle [ad_x^{-1}φ_x^{-1}p, p], e_i \rangle = 0, i = 1, . . . , r$, i.e.,

$$pr_{ann(x)\perp} \dot{p} = - pr_{ann(x)[ad_x^{-1}φ_x^{-1}p, p]} \tag{15}$$

The relations (12) and (15) proves (8).

Now, for the normal metric $K_0$ we have $φ_x = −ad_x^{-1} \circ ad_x^{-1}$ and the Hamiltonian is

$$H_0 = -\frac{1}{2}⟨ad_xadm_x, p⟩ = \frac{1}{2}⟨[x, p], [x, p]⟩ = \frac{1}{2}⟨Φ_0(x, p), Φ_0(x, p)⟩. \tag{16}$$

The equation (9) follows directly from (7), while (12) and (15) become

$$pr_{ann(x)\perp} \dot{p} = ad_x^{-1}[p, [x, [p, x]]] + ε[x, p],$$

$$pr_{ann(x)\perp} \dot{p} = pr_{ann(x)[[x, p]], p}.$$ Again, the Jacobi identity gives

$$[p, [x, [p, x]]] = [x, [[x, p], p]] = ad_x(pr_{ann(x)[[x, p]], p})$$

which together with the above formulae proves (10). \[\Box\]

The geometry of the Hamiltonian flows on cotangent bundles, in this representation, is studied by Bloch, Brockett and Crouch [1]. The system (9), (10),
for $\epsilon = 0$, agrees with the equations (2.7) given in [1], while the system \( \mathcal{F}_1 \), \( \mathcal{F}_2 \) differs from the equations (2.19) \[1\]. The equations (2.19) \[1\] describe the geodesic flows of submersion (or collective) metrics on the orbit $\mathcal{O}(a)$, and, in general, are not $G$-invariant. Recall that the submersion metrics are given by \( \mathcal{H} = \frac{1}{2} \langle \Phi_0(x, p), \phi \Phi_0(x, p) \rangle \), where $\phi$ is a symmetric, positive definite operator on $\mathfrak{g}$. Specially, $K_0$ is both $G$-invariant and submersion metric.

### 3 Integrable Flows

Let $\mathcal{F}_1^\epsilon$ be the algebra of all analytic, polynomial in momenta, functions of the form $\mathcal{F}_1^\epsilon = \{ p \circ \Phi_\epsilon | p \in \mathbb{R}[\mathfrak{g}] \}$ and $\mathcal{F}_2$ be the algebra of all analytic, polynomial in momenta, $G$-invariant functions on $T^* \mathcal{O}(a)$. Then, according to the Noether theorem

$$\{ \mathcal{F}_1^\epsilon, \mathcal{F}_2 \}^\epsilon = 0,$$

where $\{ \cdot, \cdot \}^\epsilon$ are magnetic Poisson bracket with respect to $\omega + \epsilon \rho^* \Omega$.

Consider the Hamiltonian $H_\epsilon = \frac{1}{2} \langle \Phi_\epsilon, \Phi_\epsilon \rangle \in \mathcal{F}_1^\epsilon$. A simple calculation shows $H_\epsilon(x, p) = H_0 + \epsilon^2 \frac{1}{2} \langle a, a \rangle$. Thus, we see that Hamiltonian flows of $H_0$ and $H_\epsilon$ coincide. Since $H_\epsilon$ belongs to $\mathcal{F}_1^\epsilon$ its commutes with $\mathcal{F}_2$. On the other side, as a composition of the momentum mapping with an invariant polynomial, the function $H_\epsilon$ is also $G$-invariant and commutes with $\mathcal{F}_1^\epsilon$. From the above consideration and Theorem 2.1 \[3\] we recover the Efimov result \[3\]:

**Theorem 2** Let $G$ be a compact Lie group and $a \in \mathfrak{g}$. The magnetic geodesic flows of normal metric \[9\], \[10\] on the adjoint orbit $\mathcal{O}(a)$ is completely integrable in the non-commutative sense.

Namely, the algebra of first integrals $\mathcal{F}_1^\epsilon + \mathcal{F}_2$ is complete on $(T^* \mathcal{O}(a), \omega + \epsilon \rho^* \Omega)$ (see \[3\]) and its invariant level sets are isotropic tori. Similarly as in the Liouville theorem, the tori are filled up with quasi-periodic trajectories of the system \[9\], \[10\] (see \[3\] \[11\]).

**Integrable Deformations.** Let $\mathcal{A} \subset \mathbb{R}(\mathfrak{g})$ be a commutative set of polynomials with respect to Lie-Poisson brackets on $\mathfrak{g}$. One can always find $\mathcal{A}$ that is complete on generic orbits $\mathcal{O}(\Phi_\epsilon(x, p))$ (e.g, see \[2\]). Let $\Phi_\epsilon^* \mathcal{A}$ be the pull-back of $\mathcal{A}$ by the momentum map: $\Phi_\epsilon^* \mathcal{A} = \{ h \circ \Phi_\epsilon | h \in \mathcal{A} \}$.

Let $\mathcal{B}$ be a commutative subset of $\mathcal{F}_2$, with respect to the magnetic Poisson bracket. Then $\Phi_\epsilon^* \mathcal{A} + \mathcal{B}$ is a complete commutative set on $(T^* \mathcal{O}(a), \omega + \epsilon \rho^* \Omega)$ if $\mathcal{B}$ is a complete commutative subset of $\mathcal{F}_2$, i.e., we have

$$\delta = \dim \mathcal{O}(a) - \frac{1}{2} \dim \mathcal{O}(\Phi_\epsilon(x, p))$$

independent functions in $\mathcal{B}$, for a generic element $(x, p) \in T^* \mathcal{O}(a)$ \[4\].
The $G$-invariant, polynomial in momenta functions $f(x, p)$ on $T^*O(a)$, are in one-to-one correspondence with $\text{Ad}_{G_a}$-invariant polynomials on $\mathfrak{v}$ via restriction to $T^*_aO(a)$: $f_0(p_0) = f(a, p_0)$. Next, we apply the transformation

$$f_0 \mapsto \tilde{f}, \quad f_0 = \tilde{f} \circ \Phi_0|_{x=a} = \tilde{f} \circ \text{ad}_a.$$ 

Within these identifications, from (2), (4) and Thimm’s formula for $\epsilon = 0$ [12], the magnetic Poisson bracket $\{f, g\}_\mathfrak{v}(x, p)$ corresponds to the following bracket (our notation is slightly different from Efimov’s [4])

$$\{\tilde{f}(\mu), \tilde{g}(\mu)\}_\mathfrak{v} = -\langle \mu + \epsilon a, [\nabla \tilde{f}(\mu), \nabla \tilde{g}(\mu)] \rangle,$$

where $\mu = [a, p_0], x = \text{Ad}_a p, \ p = \text{Ad}_p p_0$.

Note that $\{(\cdot, \cdot)_\mathfrak{v}\}, \lambda \in \mathbb{R}$ is a pencil of the compatible Poisson brackets on the algebra of $\text{Ad}_{G_a}$-invariant polynomials $\mathbb{R}[\mathfrak{v}]^{G_a}$. By the use of this pencil and the completeness criterion derived in [2], it is proved that the family of polynomials

$$\mathcal{B}_a = \{p_\lambda^G(\mu) = p(\mu + \lambda a), \lambda \in \mathbb{R}, \ p \in \mathbb{R}[\mathfrak{g}]^G, \ \eta \in \mathfrak{v}\}$$

is a complete commutative subset of $\mathbb{R}[\mathfrak{v}]^{G_a}$ with respect to the canonical brackets $\{(\cdot, \cdot)\}_\mathfrak{v}$ (see [3, 7]). Here $\mathbb{R}[\mathfrak{g}]^G$ is the algebra of $\text{Ad}_{G_a}$-invariant polynomials on $\mathfrak{g}$. Using the method of [2], it can be verified that $\mathcal{B}_a$ is a complete commutative set with respect to the magnetic Poisson bracket [13] as well.

Let $b$ an element from the center of $\text{ann}(a)$. Define the sectional operator $\tilde{\phi}_{a,b} : \mathfrak{v} \to \mathfrak{v}$ by $\tilde{\phi}_{a,b} = \text{ad}_a^{-1} \circ \text{ad}_b = \text{ad}_b \circ \text{ad}_a^{-1}$. For compact groups, among sectional operators we can take positive definite ones. It easily follows from [8] that the function $H_{a,b} = \frac{1}{2}(\tilde{\phi}_{a,b}(\mu), \mu)$ belongs to $\mathcal{B}_a$. The corresponding $G$-invariant function is

$$H_{a,b}(x, p) = \frac{1}{2}(\text{ad}_{b_x} p, \text{ad}_x p) = -\frac{1}{2}(\text{ad}_x \text{ad}_{b_x} p, p) = \frac{1}{2}(\phi_{x,b_x} p, p),$$

where $b_x = \text{Ad}_b x, x = \text{Ad}_a b$ and $\phi_{x,b_x} = -\text{ad}_x \text{ad}_{b_x}$. (Recall that $b$ belongs to the center of $\text{ann}(a)$ and since $G$ is compact connected Lie group, $G_a$ is also connected, so $b_x$ is well defined.) This is a Hamiltonian function of the $G$-invariant metric $K_{a,b}$:

$$K_{a,b}(\eta_1, \eta_2) = \langle (\text{ad}_{b_x})^{-1} \eta_1, \text{ad}_x^{-1} \eta_2 \rangle,$$

where $\eta_1, \eta_2 \in T_xO(a)$. Whence, we get the following statement

**Theorem 3** The magnetic geodesic flows of the metrics $K_{a,b}$ with respect to the magnetic term $\epsilon \Omega$:

$$\dot{x} = -\text{ad}_x \text{ad}_{b_x} p = [[b_x, p], x],$$

$$\dot{p} = -\text{ad}_x^{-1}[p, [x, [b_x, p]]] + \text{pr}_{\text{ann}(x)}[[b_x, p], p] + \epsilon [b_x, p]$$

are completely integrable in the commutative sense, by means of analytic, polynomial in momenta first integrals.
The Liouville Lagrangian tori are additionally foliated by $\delta$-dimensional invariant isotropic tori, level sets of integrals $F_1 + B_\eta$ ($\delta$ is given by (17)). Note that $\delta$ does not depend on $\epsilon$: for a generic $\eta \in \nu$ we have equality $\dim \mathcal{O}(\eta + \epsilon a) = \dim \mathcal{O}(\eta)$ for all $\epsilon \in \mathbb{R}$ (see [4, 7]). Therefore, the influence of the magnetic fields $\epsilon \Omega$, $\epsilon \in \mathbb{R}$ reflects as a deformation of the foliation of the phase space $T^* \mathcal{O}(a)$ by invariant tori. As the magnetic field increases, the magnetic geodesic lines become more curved.

**Concluding Remarks.** One can take $b$ such that the operator $\phi_{x,b}$ is positive, but with kernel different from zero. Then the Hamiltonian flow of $H_{a,b}$, for $\epsilon = 0$, represents an integrable sub-Riemannian geodesic flow on the orbit $\mathcal{O}(a)$ with the constraint distribution $D$ at the point $x$ given by the image

$$D_x = \phi_{x,b}(T^*_x \mathcal{O}(a)) = \text{ad}_b(\text{ann}(x))^\perp \subset T_x \mathcal{O}(a)$$

and the sub-Riemannian structure defined by (20), where now $\eta_1, \eta_2 \in D_x$. Here we assume that the distribution $D$ is bracket generating (see [10] for more details).

There is a natural generalization of the above results to the class of magnetic potential systems on coadjoint orbits as well as to the wider class of homogeneous spaces. We shall consider these problems in the forthcoming paper.

**Acknowledgments.** The first author was supported by Russian Found for Basic Research, RFBR 05-01-00978. The second author was supported by the Serbian Ministry of Science, Project ”Geometry and Topology of Manifolds and Integrable Dynamical Systems”.

**References**

[1] Bloch, A. M., Brockett, R. W. and Crouch P. E.: Double Bracket Equations and Geodesic Flows on Symmetric Spaces. Commun. Math. Phys. 187, 357-373 (1997)

[2] Bolsinov, A. V.: Compatible Poisson brackets on Lie algebras and the completeness of families of functions in involution, Izv. Acad. Nauk SSSR, Ser. matem. 55, no.1, 68-92 (1991) (Russian); English translation: Math. USSR-Izv. 38, no.1, 69-90 (1992)

[3] Bolsinov, A. V. and Jovanović, B.: Non-commutative integrability, moment map and geodesic flows. Annals of Global Analysis and Geometry 23, no. 4, 305-322 (2003), arXiv: math-ph/0109031

[4] Bolsinov, A. V. and Jovanović, B.: Complete involutive algebras of functions on cotangent bundles of homogeneous spaces. Mathematische Zeitschrift 246 no. 1-2, 213-236 (2004)

[5] Efimov, D. I. The magnetic geodesic flows in a homogeneous field on the complex projective space. Sib. Mat. Zh. 45, no. 3, 566-576 (2004) (Russian); English translation: Siberian Mathematical Journal 45, no.3, 465-474 (2004)
[6] Efimov, D. I. The magnetic geodesic flows on a homogeneous symplectic manifold. Sib. Mat. Zh. 46, no. 1, 106-118 (2005) (Russian); English translation: Siberian Mathematical Journal 46, no.1, 83-93 (2005)

[7] Mykytyuk, I. V. and Panasyuk A.: Bi-Poisson structures and integrability of geodesic flows on homogeneous spaces. Transformation Groups 9, no. 3, 289-308 (2004)

[8] Mishchenko, A. S. and Fomenko, A. T.: Euler equations on finite-dimensional Lie groups. Izv. Acad. Nauk SSSR, Ser. matem. 42, no.2, 396-415 (1978) (Russian); English translation: Math. USSR-Izv. 12, no.2, 371-389 (1978)

[9] Mishchenko, A. S. and Fomenko, A. T.: Generalized Liouville method of integration of Hamiltonian systems. Funkts. Anal. Prilozh. 12, No.2, 46-56 (1978) (Russian); English translation: Funct. Anal. Appl. 12, 113-121 (1978)

[10] Montgomery R.: A Tour of Subriemannian Geometries, Their Geodesics and Applications, Mathemtical Surveys and Monographs 91, AMS, 2002

[11] Nekhoroshev, N. N.: Action-angle variables and their generalization. Tr. Mosk. Mat. O.-va. 26, 181-198, (1972) (Russian); English translation: Trans. Mosc. Math. Soc. 26, 180-198 (1972)

[12] Thimm A.: Integrable geodesic flows on homogeneous spaces. Ergod. Th. & Dynam. Sys. 1, 495-517 (1981)