GALAXY CLUSTER CORRELATION FUNCTION TO Z ~ 1.5 IN THE IRAC SHALLOW CLUSTER SURVEY

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ABSTRACT

We present the galaxy cluster autocorrelation function of 277 galaxy cluster candidates with 0.25 ≤ z ≤ 1.5 in a 7 deg² area of the IRAC Shallow Cluster Survey. We find strong clustering throughout our galaxy cluster sample, as expected for these massive structures. Specifically, at ⟨z⟩ = 0.5 we find a correlation length of r₀ = 17.40±3.30 h⁻¹ Mpc, in excellent agreement with the Las Campanas Distant Cluster Survey, the only other non–local measurement. At higher redshift, ⟨z⟩ = 1, we find that strong clustering persists, with a correlation length of r₀ = 19.14±5.65 h⁻¹ Mpc. A comparison with high resolution cosmological simulations indicates these are clusters with halo masses of ~ 10¹⁴M☉, a result supported by estimates of dynamical mass for a subset of the sample. In a stable clustering picture, these clusters will evolve into massive (10¹⁵M☉) clusters by the present day.

Subject headings: galaxies: clusters: general — cosmology: observations — large–scale structure of the universe

1. INTRODUCTION

The clustering amplitude of massive galaxy clusters, and in particular its dependence on richness or cluster mass, is a strong function of the underlying cosmology. Recent theoretical work (e.g., Majumdar and Mohr 2003; Younger et al. 2006) has demonstrated how cluster surveys can be “self–calibrated,” providing precise simultaneous constraints on both the cosmology and cluster evolution models. Given redshift and approximate mass information for each cluster, reliable cosmological parameter estimation is feasible even in the presence of significant, and potentially unknown, evolution in cluster physical parameters (e.g., Gladders et al. 2007). The ultimate goal of these cosmological studies, a measurement of the equation of state of dark energy with an accuracy competitive with SNe Ia methods, also requires knowledge of the cluster power spectrum or autocorrelation function (Majumdar and Mohr 2004; Wang et al. 2004). Such extensive sample characterization naturally emerges from mid–infrared selected photometric redshift cluster surveys like the IRAC Shallow Cluster Survey (ISCS; Eisenhardt et al. 2007, hereafter E07).

Previous studies (e.g., Bahcall et al. 2003, and references therein) have produced largely local (z ≤ 0.3) galaxy cluster clustering measurements with limited baselines for evolutionary studies. The ISCS provides the opportunity to address these issues in one of the largest statistical samples of high redshift clusters to date. In this Letter we present the first measurement of the galaxy cluster autocorrelation function extending over more than half the cosmic age of the Universe. Such measurements also allow evolutionary connections between galaxy clusters and high redshift, highly clustered galaxy populations to be explored.

We use a concordance cosmology throughout, with ΩM = 0.3, ΩΛ = 0.7, and H₀ = 70 km s⁻¹ Mpc⁻¹. For consistency with previous studies we report distances, including correlation lengths, in units of comoving h⁻¹ Mpc, with H₀ = 100 h km s⁻¹ Mpc⁻¹.

2. IRAC SHALLOW CLUSTER SURVEY

The ISCS is a sample of 335 galaxy clusters spanning 0.1 ≤ z ≤ 2 in the IRAC Shallow Survey (ISS, Eisenhardt et al. 2004) area of the NOAO Deep–Wide Field Survey (NDWFS, Jannuzi and Dey 1999; Jannuzi et al. in prep) in Boötes. The clusters are selected using a wavelet detection algorithm which identifies peaks in cluster probability density maps constructed from accurate photometric redshift probability functions for 175,431 galaxies brighter than 13.3 µJy at 4.5 µm in a 7.25 deg² region (Brodwin et al. 2006, hereafter B06; E07). The galaxy photometric redshifts, which are key to the cluster–finding algorithm, are derived from the joint ISS, NDWFS (DR3) and FLAMEX (Elston et al. 2006) data sets.

The AGES survey (Kochanek et al. in prep) in Boötes provides spectroscopic confirmation for dozens of clusters at z ≤ 0.5 (E07). At higher redshift, a multi–year Keck spectroscopic campaign has to date confirmed 10 z > 1 clusters (Stanford et al. 2005; Elston et al. 2006; B06; E07). Extensive masking, described in B06, is implemented to reject areas suffering from remnant cosmetic artifacts or data quality issues which could affect the robustness of the clustering measurements. The final unmasked area is 7.00 deg² and contains 320 of the 335 galaxy clusters in the full ISCS, of which 277 are within the redshift range (0.25 ≤ z ≤ 1.5) considered in this Letter.

3. GALAXY CLUSTER AUTOCORRELATION MEASUREMENTS

3.1. Angular Correlation Function in Redshift Slices

The redshift distribution of galaxy clusters is presented in Figure 1. A fit of the form N(z) ∝ (z/z₀)α exp(-(z/z₀)β), with α = 1.5, β = 1.7 and z₀ = 0.7, is overplotted. Cosmological studies with this observed redshift distribution must await a detailed characterization of the survey selection function, to be presented in a forthcoming paper. Galaxy clusters are split into two similarly populated redshift bins, 0.25 ≤ z ≤ 0.75 (136 clusters) and 0.75 < z ≤ 1.5 (141 clusters), using the best available redshift information (i.e., spectroscopic redshifts where available, photometric redshifts otherwise). These bins
span the range where the accuracy, \( \sigma_z = 0.06(1 + z) \), and reliability of the galaxy photometric redshifts have been demonstrated (B06; see also Brown et al. 2007).

\[
\text{Fig. 1.} \quad \text{Observed redshift distribution of galaxy clusters in the ISCS. The dashed lines illustrate the two redshift bins considered here, } 0.25 \leq z \leq 0.75 \text{ and } 0.75 < z < 1.5, \text{ containing 277 clusters between them. The curve is a fit to the distribution and is used in the Limber deprojection. This distribution should not be used for quantifying number density evolution without careful inclusion of selection biases.}
\]

The angular correlation function (ACF) is parametrized here as a simple power law,

\[
\omega(\theta) = A_\omega \theta^{-\delta}.
\]

This can be deprojected (Limber 1954) to yield a measurement of the real-space correlation length, \( r_0(z) \), over the redshift range spanned by the 2-D sample,

\[
r_0^2(z_{\text{eff}}) = A_\omega \left[ \frac{H_0 H_\gamma}{c} \right] \frac{\int z N(z)^2 [x(z)]^{1-\gamma} E(z) \, dz}{\int z N(z)^2 \, dz} \left( \frac{1}{\Gamma(\gamma)} \right)^{-1},
\]

where

\[
z_{\text{eff}} = \int z N(z)^2 [x(z)]^{1-\gamma} E(z) \, dz / \int z N(z)^2 \, dz
\]

\( \gamma \equiv 1 + \delta, \ H_\gamma = \Gamma(1/2) \Gamma[(\gamma - 1)/2] / \Gamma(\gamma/2), \) \( N(z) \) is the redshift distribution, and \( E(z) \) and \( x(z) \) describe the evolution of the Hubble parameter and the comoving radial distance, respectively (e.g., Hogg 1999). Since the cluster redshift uncertainties are much smaller than the width of our redshift bins (E07), they have little impact on the redshift distribution and the modeling of the spatial correlation function.

3.2. Results

We calculate the ACF using both the Landy and Szalay (1993) and Hamilton (1993) estimators using 500,000 randoms to ensure a robust Monte Carlo integration. The results are nearly identical with both estimators and we report the results obtained with the latter. Two independent fitting techniques are applied to the data. The standard frequentist (or classical) approach is used to simultaneously fit the slope, and, through the use of the relativistic Limber equation, the correlation length, \( r_0 \). In calculating these correlation lengths we have adopted the \( N(z) \) parametrization shown in Figure 1. A Bayesian technique is also employed to directly determine the correlation lengths, marginalizing over the slope, \( \delta \), subject to the weak prior that it is in the range \( 0.2 \leq \delta \leq 1.8 \).

This method is desirable since, despite the large number of clusters, it is difficult to simultaneously constrain both amplitude and slope. Figure 2 shows both the frequentist fits and the Bayesian likelihood functions in \( r_0 \).

\[
\text{Fig. 2.} \quad \text{Angular correlation functions in redshift bins at } 0.25 \leq z \leq 0.75 \text{ (top) and } 0.75 < z < 1.5 \text{ (bottom). The error bars are estimated via bootstrap resampling. The red solid lines show the best-fit ACFs, with corresponding slopes listed in Table 1. The blue dashed lines show the fits with a fixed slope (} \delta = 1.0\text{). Marginalizing over the slope produces the likelihood functions in } r_0 \text{ presented in the insets. The 68% and 90% confidence intervals are shown.}
\]

In this highly clustered population, sample (or cosmic) variance dominates over simple Poisson errors. We therefore calculate the error bars using 100 bootstrap resamplings with replacement. To test for possible systematics across the field we divided the field into halves, once north–south and once east–west, and for each subfield we computed the ACF. The results for all half–fields agree within 1 \( \sigma \), indicating that we are not adversely affected by an unidentified bias in the spatial selection function.

Bootstrap simulations of the cluster detection process indicate that the spurious fraction is less than 10% at all redshifts (E07), and spectroscopic observations indicate it is likely much lower. Conservatively assuming 5–10% of the sample is indeed spurious, and that these are uncorrelated, then at most we are underestimating the clustering by \( \approx 11–23\%\).

In this work we adopt the \( r_0 \) values from the Bayesian fits,
though the fits from both methods, presented in Table 1, are completely consistent. The correlation amplitude in our low redshift bin at $z \approx 0.5$ is in excellent agreement with the only other measurement at this redshift (Gonzalez et al. 2002). Our high redshift measurement, at $z \approx 1$, is the first to probe structure on the largest scales in the first half of Universe.

The space densities for these samples and the mean intercluster distances, $d_c$, are also presented in Table 1. The relationship between clustering amplitude and $d_c$ predicted in a concordance cosmology, and observed in practice (Bahcall et al. 2003, and references therein), is only weakly dependent on redshift. As shown in Figure 3, the ISCS samples are quite consistent with the LCDM predictions between $0 < z < 1.5$ (hashed region) from Younger et al. (2005).

4. DISCUSSION

A key theoretically predictable cluster observable is the correlation function as a function of halo mass. In simulations the halo mass, $M_{200}$, is defined as the mass inside the radius at which the mean overdensity is 200 times the critical density. We compare our clustering results with the Younger et al. (2005) analysis of the Hopkins et al. (2005) high–resolution cosmological simulation, which had a 1500 $h^{-1}$ Mpc box length, an individual particle mass of $1.8 \times 10^{11} M_{\odot}$, and a power spectrum normalization of $\sigma_8 = 0.84$. We infer that the ISCS cluster sample has average log [$M_{200}/M_{\odot}$] masses of $\sim 13.9^{+0.3}_{-0.3}$ and $\sim 13.8^{+0.2}_{-0.3}$ at $z_{\text{eff}} = 0.53$ and 0.97, respectively. Direct dynamical masses for the 10 $z > 1$ clusters presented in E07 yield a largely consistent distribution of masses (Gonzalez et al. in prep).

The observed constancy of $r_0$ for massive galaxy clusters out to $z = 1$ is a robust confirmation of a key prediction from numerical simulations, reflecting the relative constancy of the mass hierarchy of clusters with redshift (Younger et al. 2005). That is, the N most massive clusters at one epoch roughly correspond to the N most massive clusters at a later epoch, and therefore have similar clustering.

In Figure 4 we plot a compilation of recent clustering amplitudes for various cluster surveys, as well as for highly clustered galaxy populations, including FRII radio galaxies (Overzier et al. 2003), EROs (Brown et al. 2005; Daddi et al. 2001, 2004), ULIRGs (Farrah et al. 2006; Magliocchetti et al. 2007), and SMGs (Blain et al. 2004). Since optically–selected QSOs are more modestly clustered (Coil et al. 2007; Croom et al. 2005; Myers et al. 2006; Porciani et al. 2004), and therefore reside in considerably less massive halos ($\sim 10^{12} - 10^{13} M_{\odot}$) than the ISCS clusters, they are not included in Figure 4.

Following Moustakas and Somerville (2002), we overplot the halo conserving model of Fry (1996) normalized to our two measurements in order to explore possible evolutionary connections with structures at other redshifts. The shaded area (dotted lines) shows the 1 $\sigma$ region for the $z_{\text{eff}} = 0.53$ (0.97) measurement. In this model, representative of a class of merger-free biased structure formation models, the ISCS clusters will evolve into typical present–day massive clusters, such as those in the SDSS, APM or Abell surveys. In the stable–clustering picture, in which clustering is fixed in physical coordinates (Groth and Peebles 1977, dashed lines), the $z_{\text{eff}} = 0.97$ ISCS clusters grow into the most massive clusters in the local Universe, typically identified in X–ray surveys.

Most of the plotted high redshift galaxy clustering measurements are rather uncertain due to both small number statistics and poorly known redshift distributions. Mindful of this caveat, we observe that FRII radio galaxies, some ERO samples, and $z \sim 2$ ULIRGs have clustering consistent with that seen in the ISCS clusters in either model. Clearly these populations trace very massive halos ($\gtrsim 10^{13} M_{\odot}$). As a caution against overinterpretation, however, we note that only
ULIRGs have space densities similar to the present cluster samples, a prerequisite for drawing evolutionary connections from these particular models. Thus the present work offers a measure of support for recent studies (e.g., Magliocchetti et al. 2007) indicating that ULIRGs may be associated with, or progenitors of, groups or low-mass clusters.

5. CONCLUSIONS

By deprojecting the angular correlation function measured in redshift bins spanning \( z = 0.25 \) to \( z = 1.5 \), we have determined the real-space clustering amplitudes for ISCS clusters at \( \bar{z} = 0.53 \) and \( \bar{z} = 0.97 \) to be \( r_0 = 17.40^{+3.16}_{-3.10} \) and \( r_0 = 19.14^{+5.56}_{-5.36} \) \( h^{-1} \) Mpc, respectively. These measurements are consistent with the relation between correlation amplitude and mean intercluster distance predicted by LCDM. The ISCS clusters have total masses exceeding \( 10^{14} M_\odot \) and will evolve into very massive clusters by the present day.

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\( \Delta z \), \( \bar{z} \)\(^{\text{eff}} \), \( N \), \( \log A_{r^0} \), \( r_0 \), \( \delta \), \( r_0 \), \( \bar{n} \), \( \bar{d} \), \( \bar{d}_c \)

| \( \Delta z \) | \( \bar{z} \)\(^{\text{eff}} \) | \( N \) | \( \log A_{r^0} \) | \( r_0 \) | \( \delta \) | \( r_0 \) | \( \bar{n} \) | \( \bar{d} \) | \( \bar{d}_c \) |
|---|---|---|---|---|---|---|---|---|---|
| 0.25 – 0.75 | 0.53 | 136 | -1.37\(+0.31\) | 18.70\(+3.31\) | 1.00\(+0.46\) | 17.40\(+3.98\) | 31.8 \pm 2.7 | 31.6 \pm 0.9 |
| 0.75 – 1.50 | 0.97 | 141 | -1.73\(+0.52\) | 16.84\(+5.18\) | 1.21\(+0.89\) | 19.14\(+5.65\) | 8.96 \pm 0.76 | 48.1 \pm 1.4 |

\(^{a}\)Formal fits for \( r_0 \) and \( \delta \) were computed directly from the data. The error in \( A_{r^0} \) corresponds to the error in \( r_0 \) at the best-fit \( \delta \).

\(^{b}\)Best-fit parameters from frequentist fits.

\(^{c}\)Best-fit \( r_0 \) from Bayesian marginalization.