Synthesis of Fe-Doped Tetrahedrites \( \text{Cu}_{12-x}\text{Fe}_x\text{Sb}_4\text{S}_{13} \) and Characterization of Their Thermoelectric Properties

Sung-Yoon Kim, Ji-Hee Pi, Go-Eun Lee, and Il-Ho Kim*

Department of Materials Science and Engineering, Korea National University of Transportation, Chungju, 27469, Republic of Korea

Abstract: Tetrahedrite (\( \text{Cu}_{12}\text{Sb}_4\text{S}_{13} \)) has attracted attention as a p-type thermoelectric material with very low thermal conductivity induced by the anharmonic oscillation of Cu due to the lone-pair electrons of Sb. Many studies have been conducted to improve its thermoelectric performance by partially substituting the transition elements for the Cu sites. In this study, Fe-doped tetrahedrites \( \text{Cu}_{12-x}\text{Fe}_x\text{Sb}_4\text{S}_{13} \) \((x = 0.1-0.4)\) were prepared by mechanical alloying and hot pressing. The tetrahedrite phase was successfully synthesized by mechanical alloying without post-annealing and exhibited stability even without phase transition after hot pressing. Moreover, the Fe content was observed to be directly proportional to the lattice constant, which confirmed the Fe substitutions on the Cu sites. The electrical conductivity was observed to decrease with the increase in the Seebeck coefficient due to the charge compensation caused by Fe doping (electron donation). The highest power factor was 0.84 mWm\(^{-1}\)K\(^{-2}\) at 723 K for the specimen with \( x = 0.1 \); however, it decreased with an increase in Fe content. In addition, as the Fe content increased, the electronic thermal conductivity decreased. Thus, the lowest thermal conductivity value was obtained for the specimen with \( x = 0.4 \) (0.45–0.64 Wm\(^{-1}\)K\(^{-1}\)) in the temperature range of 323–723 K. As a result, the maximum value of the dimensionless figure of merit (\( ZT = 0.80 \)) was achieved at 723 K for the specimen with \( x = 0.2 \).

(Received March 19, 2020; Accepted April 7, 2020)

Keywords: thermoelectric, tetrahedrite, mechanical alloying, hot pressing

1. INTRODUCTION

Tetrahedrite is an abundant and natural sulfosalt mineral, which is considered to be a promising p-type thermoelectric material. Consequently, numerous studies have recently been conducted on synthetic tetrahedrite (\( \text{Cu}_{12}\text{Sb}_4\text{S}_{13} \)) [1-3]. Tetrahedrite belongs to the space group \( \text{i}4\overline{3}m \), contains 58 atoms per unit cell, and its complex crystal structure results in intrinsically low thermal conductivity [1]. Considering the valences and structural bonding states of Cu, Sb, and S, the composition of the synthetic tetrahedrite can be represented as \( \text{Cu}^{I}_{10}\text{Cu}^{II}_{2}\text{Sb}_{4}\text{S}_{12}^{II} \). Structurally, the rattling of \( \text{Cu}^{I} \) atoms, which are located inside the triangular plane of the S atoms, occurs due to the interaction of these atoms with the lone-pair electrons of the Sb atoms, interfering with heat transfer. This leads to low lattice thermal conductivity in tetrahedrites [4-6]. Extensive research has been conducted to improve the thermoelectric properties of tetrahedrites, such as improving the dimensionless figure of merit (\( ZT \)) by partial substitution of \( \text{Cu}^{II} \), Sb, or \( \text{S}^{II} \) with other elements, specifically, \( \text{Cu}_{11.8}\text{Co}_{0.2}\text{Sb}_4\text{S}_{13} \), \( \text{Cu}_{11.8}\text{Ni}_{0.2}\text{Sb}_4\text{S}_{13} \), and \( \text{Cu}_{11.6}\text{Zn}_{0.4}\text{Sb}_4\text{S}_{13} \) with substituted Co, Ni, and Zn, respectively, on the Cu sites. At 723 K, they achieved \( ZT \) values of 0.94, 0.92, and 0.76, respectively [7-9]. Further, \( \text{Cu}_{12}\text{Sb}_{3.9}\text{Bi}_{0.1}\text{Sb}_4\text{S}_{13} \), where Bi is substituted on the Sb site, obtained a \( ZT \) value of 0.88 at 723 K [10].

Additional electrons are generated when another transition metal having an oxidation state of +2 or +3 is substituted for the \( \text{Cu}^{+} \) ion on the \( \text{Cu}^{II} \) site. As a result, the charge compensation reduces the carrier (hole) concentration, and the electronic thermal conductivity decreases due to the reduced carrier contribution. In addition, this substitution causes phonon scattering, which in turn leads to low lattice thermal conductivity. Consequently, the dimensionless figure of merit can be improved by doping with transition elements. Tipireddy et
tetrahedrites (Cu$_{12-x}$ and Cu$_{11.5}$ tetrahedrites, namely: Cu$_{11}$ magnetic properties, and thermoelectric properties of al. studied the high-temperature thermoelectric properties of Cu$_{12-x}$ (Zn/Fe)$_x$Sb$_4$Te$_3$ (TM = Fe, Mn, or Co). As a result, a ZT value of 0.98 and 0.6 was achieved for Cu$_{11.0}$Co$_{0.5}$Sb$_4$Te$_3$ and Cu$_{11.0}$Fe$_{0.5}$Sb$_4$Te$_3$, respectively, at 673 K. Lu et al. [2] studied the high-temperature thermoelectric properties of synthetic tetrahedrite by varying the Zn and Fe content in Cu$_{12-x}$Sb$_4$Te$_3$ (x = 1.0 for Zn and x = 0.5 for Fe). As a result, a ZT value of 0.8 was achieved at 720 K for Cu$_{11.0}$Fe$_{0.5}$Sb$_4$Te$_3$. Nasonova et al. [12] studied the Fe-doped tetrahedrites (Cu$_{12-x}$Fe$_x$Sb$_4$Te$_3$, where x = 0.0–2.0) using Mössbauer spectroscopy, and proposed an amended model for the Fe$^{2+}$ and Fe$^{3+}$ distributions in the iron-containing tetrahedrite. Weller et al. [13] measured the thermoelectric properties of Cu$_{12-x}$Fe$_x$Sb$_4$Te$_3$ (x = 1.0–1.5) to observe the n-type behavior at low temperatures.

Since the melting point of Cu is higher than the boiling point of S, the synthesis of tetrahedrite using a conventional melting method requires sophisticated heating, maintenance, and cooling processes. In addition, phase transformation and homogenization are only possible after applying long-term annealing (usually 1-2 weeks) [14, 15]. However, for powder synthesis, mechanical alloying as a solid-state process has several advantages over the melting process with regard to the elimination of certain processes, such as vaporization, phase segregation, and pulverization [16]. In our previous study [17], we determined the optimal processing conditions for the synthesis and consolidation of Cu$_{12.5}$Fe$_{0.5}$Sb$_4$Te$_3$ using mechanical alloying (MA) and hot pressing (HP) processes. A single tetrahedrite phase was obtained in a significantly shorter processing time without post-annealing as compared to the other studies. In this study, the Fe-doped tetrahedrites were prepared by the mechanical alloying and hot pressing (MA-HP) processes, and their thermoelectric properties were examined. When Fe was substituted for Cu, the carrier concentration was reduced due to the added electrons, and thus a high power factor is expected due to the changes in electrical conductivity and Seebeck coefficient. In addition, an improved ZT value is expected due to the decreased thermal conductivity caused by the reduction in lattice contribution, due to phonon scattering and the reduction in carrier contribution due to charge compensation.

2. EXPERIMENTAL PROCEDURE

Fe-doped tetrahedrites Cu$_{12-x}$Fe$_x$Sb$_4$Te$_3$ (x = 0.1–0.4) were synthesized by MA. The elemental powders of Cu (99.9% purity, < 45 μm in size, Kojundo Chemical Laboratory Co. Ltd.), Fe (99.9% purity, < 53 μm in size, Kojundo Chemical Laboratory Co. Ltd.), Sb (99.999% purity, < 150 μm in size, Kojundo Chemical Laboratory Co. Ltd.), and S (99.99% purity, < 75 μm in size, Kojundo Chemical Laboratory Co. Ltd.) were weighed to a stoichiometric ratio. Further, 20 g of the mixed powders and 400 g of stainless-steel balls (5 mm in diameter) were placed in a hardened steel jar. After evacuating the jar, Ar gas was injected into the jar and MA was performed at a rotational speed of 350 rpm for 24 h using a planetary mill (Fritsch Pulverisette5). The synthesized mechanically-alloyed powder was charged into a graphite mold having an inner diameter of 10 mm. Subsequently, the HP process was carried out in a vacuum at a temperature of 723 K and a pressure of 70 MPa for 2 h. Detailed conditions of the MA-HP processes have been presented in our previous study [17].

Phases of the HP specimens were confirmed using X-ray diffraction (XRD; Bruker D8-Advance) with Cu Kα radiation ($\lambda = 0.15405$ nm). Diffraction patterns were recorded in the 0-2θ mode (2θ = 10–90°) with a step size of 0.02° and a scan speed of 0.4 s/step. Scanning electron microscope (SEM; FEI Quanta400) and energy dispersive spectrometer (EDS; Bruker Quantax200) were used to observe and analyze the microstructures of the sintered specimens. Elemental mapping was conducted to confirm the homogeneous distribution of the constituent elements.

The Hall coefficient, carrier concentration, and mobility of the HP specimens were measured using the van der Pauw method (Keithley 7065); a constant magnetic field of 1 T and a constant current of 100 mA was applied. The Seebeck coefficient and electrical conductivity were measured using the differential temperature method and the DC four-probe method (Ulvac-Riko ZEM-3) in a He atmosphere and at a temperature range from 323 to 723 K. The thermal conductivity was determined using thermal diffusivity, specific heat, and density, measured using the laser flash method (Ulvac-Riko TC-9000H). Finally, the power factor and dimensionless figure of merit were evaluated.
3. RESULTS AND DISCUSSION

Fig. 1 shows the results of XRD analysis according to the Fe content of the sintered Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ specimens. All specimens appeared as a single tetrahedrite phase without secondary phases. Rietveld refinement using the TOPAS program revealed that as the Fe content in Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ increased, the lattice constant increased from 1.0336 nm for $x = 0.1$ to 1.0350 nm for $x = 0.4$ (Table 1). Nasonova et al. [12] have also reported that as the Fe content in Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ increased, the lattice constant increased from 1.0320 nm for Cu$_{11.5}$Fe$_{0.5}$Sb$_4$S$_{13}$ to 1.0378 nm for Cu$_{10.5}$Fe$_{1.5}$Sb$_4$S$_{13}$. Tippireddy et al. [11] also found that the lattice constant increased from 1.0320 nm for the synthetic tetrahedrite Cu$_{12}$Sb$_4$S$_{13}$ to 1.0330 nm for the Fe-doped tetrahedrite Cu$_{11.5}$Fe$_{0.5}$Sb$_4$S$_{13}$. However, these results are contradictory to the assumption that the lattice constant decreases when smaller Fe$^{3+}$ ions (ionic radius: 49 pm) are substituted at the Cu$^{+}$ sites (ionic radius: 60 pm). Friese et al. [18] suggested that there are ten Cu$^{+}$ and two Cu$^{2+}$ ions (ion radius: 57 pm) in a tetrahedrite structure, and Fe substitution on the Cu site is more stable for the combination of Fe$^{3+}$ and Cu$^{+}$ than the combination of Fe$^{2+}$ and Cu$^{2+}$ [19]. As a result, Cu$^{2+}$ ions are induced to substitute Cu$^{+}$ ions. Consequently, Fe doping results in an increase in the number of Cu$^{+}$ ions, which have a larger ionic radius than Cu$^{2+}$ ions, which in turn results in an increase in the lattice constant.

Fig. 2 presents the SEM images of the fractured surfaces of Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$. No remarkable differences can be observed in the microstructures based on the Fe content. Highly-dense sintered bodies were obtained with the HP process. Fig. 3 shows the EDS elemental maps of Cu$_{11.8}$Fe$_{0.2}$Sb$_4$S$_{13}$ as a representative sample. All the elements were confirmed to be homogeneously distributed. As shown in Table 1, all the specimens have high relative densities greater than 98.8%, without significant differences between the actual and nominal compositions. However, the Fe content is observed to be slightly higher than the nominal composition, which is considered to be due to contamination from the steel jar and steel balls during the MA process.

![Fig. 1. XRD patterns of mechanically-alloyed and hot-pressed Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ tetrahedrites.](image1)

![Fig. 2. SEM images of fractured surfaces of Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$.](image2)

**Table 1.** Chemical compositions and physical properties of Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$.

| Composition | Relative Density [%] | Lattice Constant [nm] | Lorenz Number $[10^{-8}$ V K$^{-2}$] |
|-------------|----------------------|-----------------------|-----------------------------------|
| Cu$_{11.9}$Fe$_{0.1}$Sb$_4$S$_{13}$ | 99.0 | 1.0364 | 1.81 |
| Cu$_{11.4}$Fe$_{0.6}$Sb$_4$S$_{13}$ | 100.0 | 1.0343 | 1.77 |
| Cu$_{11.7}$Fe$_{0.3}$Sb$_4$S$_{13}$ | 100.0 | 1.0347 | 1.69 |
| Cu$_{11.6}$Fe$_{0.4}$Sb$_4$S$_{13}$ | 98.8 | 1.0350 | 1.61 |
state synthesis by MA-HP process suppressed the volatilization of constituent elements, particularly in the case of S.

The carrier concentration and mobility of Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ are shown in Fig. 4. It can be observed that the mobility tends to decrease with an increase in Fe content, whereas the carrier concentration does not change significantly with Fe content. This was inconsistent with the expectation that the additional electrons would decrease the carrier (hole) concentration following Fe substitution on the Cu site. As the Fe content increased, the mobility decreased due to the increase in charge carrier scattering, which can decrease electrical conductivity. However, the Hall coefficients, measured in this study, exhibited low reliability: the same sample was measured 20 times and extremely different values were obtained. Other researchers have also reported failure when attempting to measure the Hall coefficient of tetrahedrites, as the Hall coefficients were extremely low and close to zero, or extremely different values were obtained when measured again using the same sample [2,11].

The thermoelectric properties of Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ were measured in the temperature range from 323 to 723 K. Fig. 5 shows the temperature dependence of the Seebeck coefficient ($\alpha$) for Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$. Positive Seebeck coefficients were observed in all specimens, suggesting that the major carriers are holes of p-type semiconductors. The Seebeck coefficient of the p-type semiconductor is expressed as $\alpha = \frac{8}{3}\pi^2k_B^2m^*T^{-1}h^{-2}(\pi^2n)^{2/3}$ (where $k_B$ is the Boltzmann constant, $m^*$ is the effective carrier mass, $T$ is the absolute temperature, $e$ is the electronic charge, $h$ is the Planck constant, and $n$ is the carrier concentration) [20]. For the specimens with $x = 0.1$–0.3, the Seebeck coefficient increased with an increase in temperature. However, for $x = 0.4$, the Seebeck coefficient decreased with an increase in temperature. As the temperature increases, the Seebeck coefficient of the p-type semiconductor increases; however, if the temperature exceeds a certain temperature (an intrinsic transition temperature), the carrier concentration increases rapidly, and as a result, the Seebeck coefficient decreases.

Therefore, as the temperature of a p-type semiconductor increases, the Seebeck coefficient decreases after achieving a peak value at a specific temperature. In this study, the intrinsic transition occurred at temperatures above 723 K for $x = 0.1$–0.2 and at 623 K for $x = 0.3$. However, for $x = 0.4$, the intrinsic transition had already occurred at temperatures below 323 K. The intrinsic transition temperature is closely related to the bandgap energy, Fermi level, and carrier concentration.

At constant temperature, the Seebeck coefficient increased with an increase in the Fe content. For instance, the value increased from 179 $\mu$VK$^{-1}$ for Cu$_{11.8}$Fe$_{0.2}$Sb$_4$S$_{13}$ to 224 $\mu$VK$^{-1}$ for Cu$_{11.4}$Fe$_{0.6}$Sb$_4$S$_{13}$ at 723 K. In this study, Cu$_{11.4}$Fe$_{0.6}$Sb$_4$S$_{13}$ exhibited the highest Seebeck coefficient values of 258–224 $\mu$VK$^{-1}$ over a temperature range of 323 to 723 K. The increase in the Seebeck coefficients of tetrahedrites, caused by the transition elements substituted on the Cu sites, have been previously reported. Tipireddy et al. [11] obtained the
Seebuck coefficient value of 130 μVK⁻¹ for an undoped tetrahedrite Cu₁₂SbS₁₃ at 673 K. However, after doping with various transition elements (Mn, Fe, Co, Ni, and Zn), higher Seebuck coefficient values were obtained, such as 170 μVK⁻¹ for Cu₁₁.₅Fe₀.₅SbS₁₃ and 215 μVK⁻¹ for Cu₁₁ZnSbS₁₃ at 673 K. Heo et al. [21] also found that the Seebeck coefficient was as low as 74 μVK⁻¹ for the undoped tetrahedrite Cu₁₂SbS₁₃ at 325 K, but it increased to 170 μVK⁻¹ for Cu₁₀Fe₂SbS₁₃ and 253 μVK⁻¹ for Cu₁₀Zn₂SbS₁₃ at 325 K. These high Seebeck coefficients were attributed to the decrease in the carrier (hole) concentration, because the divalent or trivalent transition elements were substituted on the Cu⁺ sites.

Figure 6 presents the temperature dependence of the electrical conductivity (σ) in Cu₁₂₋ₓFeₓSbS₁₃. The electrical conductivity increased with an increase in temperature, and its temperature dependence significantly increased as the Fe content increased. This indicated that the Fe-doped tetrahedrites behaved as non-degenerate semiconductors. Tippireddy et al. [11] reported that the electrical resistivity of Cu₁₂₋ₓFeₓSbS₁₃ decreased (i.e., the electrical conductivity increased) with an increase in temperature. Weller et al. [13] also discovered similar results for Cu₁₂₋ₓFeₓSbS₁₃ (x = 1.0, 1.3, and 1.5). The electrical conductivity at a constant temperature decreased with an increase in Fe content, which can be attributed to the effect of charge compensation due to Fe doping [17,21].

In this study, the highest electrical conductivity value of (1.9–2.6) × 10⁴ Sm⁻¹ was obtained for Cu₁₁.₅Fe₀.₅SbS₁₃, while the lowest electrical conductivity value of (1.9–9.9) × 10³ Sm⁻¹ was obtained for Cu₁₁ₓFe₀.₁SbS₁₃, at temperatures ranging from 323 to 723 K.

Figure 7 shows the temperature dependence of the power factor for Cu₁₂₋ₓFeₓSbS₁₃. The power factor (PF) is expressed as PF = α²σ and increases proportionately with the Seebeck coefficient and the electrical conductivity [22]. Since the carrier concentration affects both Seebeck coefficient (inverse-proportionally) and electrical conductivity (proportionally), optimization of the carrier concentration is necessary to obtain high PF values. In this study, the PF increased with an increase in temperature in the measured temperature range, and it decreased as the Fe content increased. The decrease in electrical conductivity due to Fe doping was dominant over the increase in the Seebeck coefficient. Although Cu₁₁.₅Fe₀.₅SbS₁₃ exhibited the lowest Seebeck coefficient values, it exhibited the highest PF values of 0.35–0.84 mWm⁻¹K⁻² in the temperature range of 323 to 723 K, as it exhibited the highest electrical conductivity values. Tippireddy et al. [11] reported a maximum PF of approximately 1.0 mWm⁻¹K⁻² for Cu₁₁.₅Fe₀.₅SbS₁₃ at 673 K.

Figure 8 presents the temperature dependence of the thermal conductivity (κ) for Cu₁₂₋ₓFeₓSbS₁₃. The lattice thermal conductivity (κₗ) from phonon contributions was obtained by subtracting the electronic thermal conductivity (κₑ) of carrier contributions from the total thermal conductivity. Here, κₑ was estimated using the Wiedemann-Franz law (κₑ = LσT, where L is the Lorenz number) [23]. The Lorenz number [10⁻⁸
V²K⁻²] was calculated using the formula \( L = 1.5 + \exp(-|\alpha|/116) \). As shown in Table 1, as the Fe content increased, the Lorenz number decreased from \( 1.81 \times 10^{-8} \) to \( 1.61 \times 10^{-8} \) V²K⁻² at room temperature. The electronic thermal conductivity increased with an increase in temperature. However, the total thermal conductivity increased with an increase in temperature and then decreased after obtaining a peak value at 623 K, except for Cu₁₁.₆Fe₀.₄Sb₄S₁₃. This was due to the opposition between the increase in electronic thermal conductivity and the minimal changes in the lattice thermal conductivity with an increase in temperature up to 623 K. The electronic thermal conductivity was directly correlated with the electrical conductivity (carrier concentration) and the Lorenz number. Thus, it decreased with increasing Fe content.

In the case of lattice thermal conductivity, Cu₁₁.₆Fe₀.₄Sb₄S₁₃ exhibited lower values at temperatures below 523 K due to the enhanced ionized impurity scatterings caused by the dopants. Therefore, Cu₁₁.₆Fe₀.₄Sb₄S₁₃ had the lowest total thermal conductivity value of 0.45–0.64 Wm⁻¹K⁻¹ for temperatures ranging from 323 to 723 K. Cu₁₁.₆Fe₀.₄Sb₄S₁₃, which was studied by Tippireddy et al. [11] exhibited \( \kappa = 0.80–1.10 \) Wm⁻¹K⁻¹ in the temperature range of 323 to 623 K, with \( k_L = 0.60–0.75 \) Wm⁻¹K⁻¹. In this study, the fine tetrahedrite powders and the absence of the secondary phases in the specimens processed by MA were interpreted to be the reasons for the relatively low values of thermal conductivity.

The dimensionless figure of merit (\( ZT = \alpha^2 \sigma / \kappa T \)) values obtained for Cu₁₂-xFeₓSb₄S₁₃ are shown in Fig. 9. As the temperature increased, the ZT values increased. This was attributed to the increase in the power factor and the maintenance of low thermal conductivity. As the Fe content increased, both the PF and thermal conductivity decreased, resulting in the highest ZT value of 0.80 at 723 K for Cu₁₁.₈Fe₀.₂Sb₄S₁₃. Pi et al. [25] reported a ZT value of 0.66 at 723 K for the undoped Cu₁₂Sb₄S₁₃ prepared by mechanical alloying and hot pressing with the same processing conditions used in this study. Tippireddy et al. [11] achieved a ZT value of 0.6 at 673 K for Cu₁₁.₆Fe₀.₄Sb₄S₁₃. Lu et al. [26] obtained a ZT value of 0.8 at 700 K for Cu₁₁.₆Fe₀.₄Sb₄S₁₃, synthesized by encapsulated melting and
hot pressing process. Tippireddy et al. [11] utilized several days to prepare the tetrahedrite by melting process; initially, it was maintained at 973 K for 3 h, then cooled to 823 K for 30 h, and then further cooled to room temperature, followed by subsequent annealing at 773 K for 30 h. Lu et al. [26] also utilized a very slow and sophisticated melting process, of heating up to 1023 K at a rate of 0.3 K/min, holding for 12 h and then cooling at a rate of 0.5 K/min, followed by subsequent annealing at 723 K for 1 week. However, in this study, the Fe-doped tetrahedrite was successfully synthesized using the MA process for 24 h without subsequent annealing, and it exhibited superior ZT values. Therefore, the MA-HP processes were confirmed to be fast and facile methods of preparing tetrahedrite compounds with high thermoelectric performance.

4. CONCLUSION

In this study, Fe-doped tetrahedrites Cu$_{12-x}$Fe$_x$Sb$_4$S$_{13}$ ($x = 0.1, 0.2, 0.3,$ and $0.4$) were synthesized by mechanical alloying and sintered by hot pressing. The tetrahedrite compounds comprising relative densities higher than 98.8% were produced without subsequent annealing. The positive values of the Hall and Seebeck coefficients indicated that all the specimens were p-type semiconductors. The electrical conductivity exhibited positive temperature dependence, and its temperature dependence significantly increased as the Fe content increased. The Seebeck coefficient also increased with an increase in temperature, except for Cu$_{11.6}$Fe$_{0.4}$Sb$_4$S$_{13}$. At a given temperature, as the Fe content was increased, more electrons were supplied, which resulted in the decrease in electrical conductivity and the increase in Seebeck coefficient due to charge compensation. The electronic thermal conductivity also decreased in accordance with the trend in electrical conductivity. The total thermal conductivity was also reduced, exhibiting the lowest value for Cu$_{11.6}$Fe$_{0.4}$Sb$_4$S$_{13}$. The power factor decreased with increasing Fe content due to the reduction in electrical conductivity. As a result, a maximum ZT value of 0.80 was achieved at 723 K for Cu$_{11.8}$Fe$_{0.2}$Sb$_4$S$_{13}$. The MA-HP processes were demonstrated to be suitable for fabricating doped tetrahedrites with excellent thermoelectric properties.

ACKNOWLEDGMENTS

This study was supported by the Industrial Core Technology Development Program funded by the Ministry of Trade, Industry and Energy (Grant No. 10083640), and by the Basic Science Research Capacity Enhancement Project (National Research Facilities and Equipment Center) through the Korea Basic Science Institute funded by the Ministry of Education (Grant No. 2019R1A6C1010047), Republic of Korea.

REFERENCES

1. K. Suekuni, K. Tsuruta, T. Ariga, and M. Koyano, Appl. Phys. Exp. 5, 1201 (2012).
2. X. Lu, D. T. Morelli, Y. Xia, F. Zhou, V. Ozolins, H. Chi, X. Zhou, and C. Uher, Adv. Energy Mater. 3, 342 (2013).
3. T. Barbier, P. Lemoine, S. Gascoin, O. I. Lebedev, A. Kaltzoglou, P. Vaqueiro, A. V. Powell, R. I. Smith, and E. Guilmeau, J. Alloys Compd. 634, 253 (2015).
4. Y. Bouyrie, C. Candolfi, S. Pailhès, M. M. Koza, B. Malaman, A. Dauscher, J. Tobola, O. Boisron, L. Saviot, and B. Lenoir, Phys. Chem. Chem. Phys. 17, 19751 (2015).
5. X. Lu, W. Lai, Y. Wang, and D. T. Morelli, Adv. Funct. Mater. 25, 3648 (2015).
6. D. T. Morelli, E. Lara-Curzio, A. F. May, O. Delaire, M. A. McGuire, X. Lu, C. Y. Liu, and E. D. Case, J. Appl. Phys. 115, 193515 (2014).
7. S. Y. Kim, G. E. Lee, and I. H. Kim, J. Korean Phys. Soc.
Sung-Yoon Kim, Ji-Hee Pi, Go-Eun Lee, and Il-Ho Kim, *J. Electron. Mater.* **49**, 2775 (2020).

8. S.-Y. Kim, G.-E. Lee, and I.-H. Kim, *J. Electron. Mater.* **49**, 2775 (2020).

9. S.-G. Kwak, G.-E. Lee, and I.-H. Kim, *Korean J. Met. Mater.* **57**, 328 (2019).

10. S.-G. Kwak, J.-H. Pi, G.-E. Lee, and I.-H. Kim, *Korean J. Met. Mater.* **58**, 272 (2020).

11. S. Tippireddy, R. Chetty, M. H. Naik, M. Jain, K. Chattopadhyay, and R. C. Mallik, *J. Phys. Chem.* **122**, 8735 (2018).

12. D. I. Nasonova, I. A. Presniakov, A. V. Sobolev, V. Y. Verchenko, A. A. Tzirlin, Z. Wei, E. V. Dikarev, and A. V. Shevelkov, *J. Sol. Sta. Chem.* **235**, 28 (2016).

13. D. P. Weller, G. E. Kunkel, A. M. Ochs, D. T. Morelli, and M. E. Anderson, *Mater. Today Phys.* **7**, 1 (2018).

14. L. L. Huang, J. Zhang, Z. M. Wang, X. G. Zhu, J. M. Li, C. Zhu, D. Li, C. J. Song, H. X. Xin, and X. Y. Qin, *Materialia* **3**, 169 (2018).

15. A. F. May, O. Delaire, J. L. Niedziela, E. Lara-Curzio, M. A. Susner, D. L. Abernathy, M. Kirkham, and M. A. McGuire, *Phys. Rev. B* **93**, 064104 (2016).

16. Y. Q. Yu, B. P. Zhang, Z. H. Ge, P. P. Shang, and Y. X. Chen, *Mater. Chem. Phys.* **131**, 1 (2011).

17. S. Y. Kim, S. G Kwak, J. H. Pi, G. E. Lee, and I. H. Kim, *J. Electron. Mater.* **48**, 1857 (2019).

18. K. Friese, A. Grzechew, E. Makovicky, T. B. Zunic, and S. K. Moller, *Phys. Chem. Miner.* **35**, 455 (2008).

19. E. Makovicky, K. Forcher, W. Lottermoser, and G. Amthauer, *Miner. Petrol.* **43**, 73 (1990).

20. G. J. Snyder and E. S. Toberer, *Nat. Mater.* **7**, 105 (2008).

21. J. Heo, G. Laurita, S. Muir, M. A. Subramanian, and D. A. Keszler, *Chem. Mater.* **26**, 2047 (2014).

22. X. Yan, B. Poudel, Y. Ma, W. Liu, G. Joshi, H. Wang, Y. Lan, D. Wang, G. Chen, and Z. Ren, *Nano Lett.* **10**, 3373 (2010).

23. H. Cailat, A. Borschchevsky, and J. P. Fleurial, *J. Appl. Phys.* **80**, 4442 (1996).

24. B. Madaval and S. J. Hong, *J. Electron. Mater.* **45**, 12 (2016).

25. J.-H. Pi, G.-E. Lee, and I.-H. Kim, *J. Electron. Mater.* **49**, 2710 (2020).

26. X. Lu, Ph. D. Thesis, pp.61-70, Michigan State University, USA (2014).