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Partial muon capture rates in $A=3$ and $A=6$ nuclei with chiral effective field theory

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Searches for neutrinoless double-$\beta$ decay rates are crucial in addressing questions within fundamental symmetries and neutrino physics. The rates of these decays depend not only on unknown parameters associated with neutrinos, but also on nuclear properties. In order to reliably extract information about the neutrino, one needs an accurate treatment of the complex many-body dynamics of the nucleus. Neutrinoless double-$\beta$ decays take place at momentum transfers on the order of 100 MeV/$c$ and require both nuclear electroweak vector and axial current matrix elements. Muon capture, a process in the same momentum transfer regime, has readily available experimental data to validate these currents. In this Letter, we present results of ab initio calculations of partial muon capture rates for $^3$He and $^6$Li nuclei using variational and Green’s function Monte Carlo computational methods. We estimate the impact of the three-nucleon interactions, the cutoffs used to regularize two-nucleon interactions, and the energy range of $2N$ scattering data used to fit these interactions.

Introduction, conclusions, and outlook. Nuclei play a crucial role in high-precision tests of the standard model and searches for physics beyond the standard model. These investigations, including neutrinoless double-$\beta$ decay ($0\nu\beta\beta$) searches [1] and high-precision $\beta$-decay experiments [2–4], require a thorough understanding of standard nuclear effects in order to separate them from new physics signals. In particular, $0\nu\beta\beta$-decay experiments aim to establish the origin and nature of neutrino masses and test leptogenesis scenarios leading to the observed matter-anti-matter asymmetry in the universe [1]. Rates of these decays depend not only on unknown neutrino parameters, but also on nuclear matrix elements. The latter can be provided only from theoretical calculations. Thus, a prerequisite of this experimental program is an accurate treatment of the complex many-body dynamics of the nucleus and its interactions with neutrinos. If one assumes that $0\nu\beta\beta$-decay proceeds via the exchange of a light Majorana neutrino between two nucleons, then the momentum carried by the neutrino is on the order of 100 MeV/$c$ [1,5]. Muon captures on nuclei—processes where a muon captures on a proton in the nucleus releasing a neutron and a neutrino—involves momentum transfer on the order of the muon mass. The scope of this Letter is to validate our nuclear model in this kinematic regime by calculating muon capture rates in $A=3$ and $A=6$ nuclei for comparison with available experimental data.

Muon capture reactions have been treated extensively from both the theoretical and the experimental points of view [6–9], and rates have been obtained in light systems with several methods [10–21]. Here, we present calculations of partial muon capture rates using quantum Monte Carlo (QMC) methods [22]—both variational (VMC) and Green’s function Monte Carlo (GFMC) methods—to solve the nuclear many-body problem. QMC methods allow one to fully retain the complexity of many-body physics and have been successfully applied to study many nuclear electroweak properties over a wide range of energy and momentum transfer, including total muon capture rates in $^1$H and $^3$He [23], low-energy electroweak transitions [24–28], nuclear responses induced by electrons and neutrinos [29–31], $0\nu\beta\beta$-decay matrix elements [32–34], and matrix elements for dark matter scattering [35].

The Norfolk two-nucleon ($2N$) and three-nucleon ($3N$) ($NV2+3$) local chiral interactions [36–39] have been used in combination with QMC methods to study static properties of light nuclei [37,40–43], and in auxiliary-field diffusion Monte Carlo [44], Brueckner-Bethe-Goldstone [45,46], and Fermi hypernetted chain/single-operator chain [47,48] approaches to investigate the equation of state of neutron matter [49,50]. Reference [51], a study which included the current authors, reports on Gamow-Teller (GT) matrix elements calculated for $A \leq 10$ nuclei using the NV2+3 models and their consistent axial-vector currents at tree level from Refs. [38,39,52]. The study validated the many-body interactions and currents in the limit of vanishing momentum transfer. In the present Letter, we use the same nuclear Hamiltonians and axial currents, along with chiral vector currents retaining loop corrections developed in Refs. [24,25,53,54] to test the model at moderate momentum transfers on the order of 100 MeV/$c$ and to assess the sensitivity of partial muon capture rates to the dynamical input.
In the $A = 3$ system, we obtain an average rate for all Nor-
folk models of $\Gamma (A = 3; \text{VMC}) = 1512 \pm 32 \text{ s}^{-1}$ at the VMC
level that agrees with the experimental result of 1496.0 $\pm$
4.0 s$^{-1}$ [55] within error bar. In the $A = 6$ system, the
partial muon capture rate of $\Gamma (A = 6; \text{VMC}) = 1243 \pm 59 \text{ s}^{-1}$
is significantly smaller than the available experimental data point of
$1600^{+130}_{-120} \text{ s}^{-1}$ [56] but falls into the range of previous theo-
etrical estimates [57–63]. We analyzed uncertainties due to
(i) the choice of cutoffs used to regularize the NV2 interac-
tions, (ii) the energy range of 2N scattering data used to fit
the model low-energy constants (LECs), (iii) two different
versions of NV3 interactions, i.e., the nonstared model fit to the
$nd$ scattering length and the trinucleon binding energies,
and the starred model fit to the triton GT matrix element and
the trinucleon binding energies, and (iv) a 10% variation in
the nucleonic axial radius. In the $A = 3$ system, the largest
source of uncertainty comes from the choice of the 3N interaction
model, whereas for $A = 6$ we find that the uncertainty due to
the 3N interaction is slightly less than but on the order of the
cutoff and energy range uncertainties. On average, there is a
change in the rate by $\pm 0.6\%$ when the axial radius is varied
in the interval $r_A = [0.5859, 0.7161] \text{ fm}$.

We improved upon our VMC estimate by performing
GFMC propagations using models NV2+3-Ia and NV2+3-
Ia*, or Ia and Ia* for short in both the $A = 3$ and the
$A = 6$ systems. These models share the same 2N interac-
tion but differ in how the 3N interaction is fit with
$(c_D, c_E) = (3.666, -1.638)$ for Ia [40] and 
$(c_D, c_E) = (-0.635, -0.090)$ for Ia* [39]. Model Ia, constrained
by strong interaction data only, achieves 1.5% agreement with
the experimental datun for $A = 3$ with a calculated rate of
1519 $\pm 3 \text{ s}^{-1}$. Its counterpart, model Ia*, constrained to both
strong and electroweak data, underpredicts the experimental
rate by a few percent. For $A = 6$, we find that the model
Ia* propagation significantly decreases the rate due to the
monotonic growth of the $^3$He ground state rms radius at early
imaginary times. By contrast, model Ia has a stable radius
throughout the GFMC propagation, and the rate decreases by
less than 1%; nevertheless, it still underpredicts the experi-
mental datun.

Given the large error bars on the $^6$Li datum and the wide
range of values from past theoretical calculations, we ad-
vocate for renewed experimental and theoretical attention to
this partial capture rate. Whereas in this Letter we fo-
cus on $^3$He and $^6$Li to demonstrate the impact of this sort
of study, there are other muon capture rates with available
experimental data which the combination of QMC methods
and NV2+3 chiral Hamiltonians could be made to address
with future development; examples are $^{10}$B, $^{11}$B, $^{12}$C, $^{15}$O,
and $^{40}$Ca [7]. Calculations of these rates, particularly, for the
heavier nuclei, would be valuable in further validating the present
ab initio approach in the kinematic regime relevant to $0\nu\beta$-decay.

Partial muon capture rate. The muon capture processes
$^3$He$^-(\mu^-, \nu_\mu)\frac{3}{2}^1\text{He}$ and $^6$Li$^-(\mu^-, \nu_\mu)\frac{3}{2}^1\text{He}$ are induced by the
weak-interaction Hamiltonian $H_{\text{W,46,65}}$
\begin{equation}
\langle k_\nu, h_\nu|H_\text{W}|k_\mu, s_\mu \rangle = \frac{G_V}{\sqrt{2}} \int d^3x \ e^{-i\hat{\kappa}_\nu\hat{A}_\nu(x)} J^\nu(x), \tag{1}
\end{equation}
where $G_V = G_F \cos \theta_C = 1.1363 \times 10^{-5} \text{ GeV}^{-2}$ is the
Fermi coupling constant extracted from analyses of super-
lowed $\beta$-decays [66], $J^\nu$ and $\hat{A}_\nu$ are the hadronic and leptonic
four-current density operators [10], $s_\mu$ is the muon spin, $h_\nu$
is the neutrino helicity, and $k_\mu$ and $k_\nu$ are the muon and neutrino
momenta, respectively. The value of $G_V$ adopted here is from
a more recent analysis and is $\approx 1.1\%$ smaller than that used in
previous calculations based on the hyperspherical harmonics
method with chiral currents from Ref. [12].

For a transition from an initial nuclear state $\mid i, J, M_i \rangle$
where $J_{i/f}$ and $M_{i/f}$ denote the nuclear spin and its
projection—to a final nuclear state $\mid f, J_f M_f, -k_\nu \rangle$
recoiling with momentum $-k_\nu$, the general expression for the capture
rate ($\Gamma$), summed over the final states and averaged over
the initial states, is given (in the limit of vanishing $k_\nu$) by
[10,64,65]
\begin{equation}
\frac{d\Gamma}{d\omega} = \frac{1}{2(2J_i + 1)} \sum_{x_\nu, M_\nu, h_\nu, M_f} 2\pi \delta(\omega)
\times |\langle k_\nu, h_\nu; f, J_f M_f, -k_\nu|H_\text{W}|s_\mu; i, J_i M_i \rangle|^2
\frac{d^3k_\nu}{(2\pi)^3}, \tag{2}
\end{equation}
where the argument of the $\delta$ function is
\begin{equation}
\omega = E_\nu + \sqrt{E_\nu^2 + (m_f + E_f)^2 - (m_\mu + m_i + E_i)}, \tag{3}
\end{equation}
and $E_i$ and $E_f$ are the initial and the final state energies of the
nucleus [67–69]—we have neglected internal electronic energies
since they are on the order of tens of eVs for the light atoms
under consideration. We also used the following de-
finitions:
\begin{equation}
m_i = Zm_p + Nm_n + (Z - 1)m_e, \tag{4}
\end{equation}
\begin{equation}
m_f = (Z - 1)(m_p + m_n) + (N + 1)m_n \tag{4}
\end{equation}
for an initial nucleus with charge number $Z$ and neutron
number $N$, and we denoted with $m_p$, $m_n$, and $m_e$ the proton,
neutron, and electron masses, respectively.

The final integrated rate can be conveniently written in
terms of matrix elements of the nuclear electroweak current
components [11],
\begin{equation}
\Gamma = \frac{G_V^2}{2\pi} \left| \frac{\nu_{\text{av}}}{\nu_{\text{av}}} \right|^2 \frac{E_{\nu_{\text{av}}}^2}{(2J_i + 1)} \text{recoil} \sum_{M_f, M_i} \left| \langle J_f, M_f | \rho(E_\nu^2, \omega) | J_i, M_i \rangle \right|^2 + \left| \langle J_f, M_f | j_z(E_\nu^2, \omega) | J_i, M_i \rangle \right|^2
+ 2 \text{Re} \left[ \langle J_f, M_f | \rho(E_\nu^2, \omega) | J_i, M_i \rangle \langle J_f, M_f | j_z(E_\nu^2, \omega) | J_i, M_i \rangle^* \right]
+ \left| \langle J_f, M_f | j_z(E_\nu^2, \omega) | J_i, M_i \rangle \right|^2 - 2 \text{Im} \left[ \langle J_f, M_f | j_z(E_\nu^2, \omega) | J_i, M_i \rangle \langle J_f, M_f | j_z(E_\nu^2, \omega) | J_i, M_i \rangle^* \right], \tag{5}
\end{equation}
where we have chosen $\mathbf{k}_e = -\mathbf{z}$ and have introduced the outgoing neutrino energy [11],

$$E_\nu^* = \frac{(m_i + E_i + m_{\mu})^2 - (m_f + E_f)^2}{2(m_i + E_i + m_{\mu})}, \quad (6)$$

and recoil factor,

$$\frac{1}{\text{recoil}} = \left(1 - \frac{E_\nu^*}{m_i + E_i + m_{\mu}}\right). \quad (7)$$

The factor $|\psi_{1s}(0)|^2$ is written as $R|\psi_{1s}(0)|^2$, where $\psi_{1s}(0)$ is the 1s wave function, evaluated at the origin of a hydrogenlike atom, and $R$ approximately accounts for the finite size of the nuclear charge distribution [11], here calculated with the NV2+3 Hamiltonians.

**Nuclear Hamiltonians and electroweak currents.** To calculate the nuclear matrix elements required by Eq. (5) we employ VMC [70] and GFMC [71] methods. For a comprehensive review of these methods, see Refs. [22,41] and references therein. Details about the calculation of matrix elements using GFMC wave functions are found in Eqs. (19)–(24) of Ref. [72].

The many-body Hamiltonian is composed of a (one-body) kinetic energy term, and the Norfolk 2N and 3N local interactions that include next-to-next-to-next-to-leading order and next-to-next-to-leading order terms in the chiral expansion, respectively. Details about the derivation of the interaction in chiral effective field theory can be found in Refs. [36–39].

Here, we briefly summarize the differences between the model classes employed in this Letter. Models in class I (II) fit the 2N interaction to about 2700 (3700) data points up to laboratory energy of 125 (200) MeV in the nucleon-nucleon scattering database with a $x^2$/datum of about $\lesssim 1.1$ ($\lesssim 1.4$). Within each class, models a and b differ in the set of cutoffs adopted to regularize the short- and long-range components of the interaction either $(R_S, R_L) = (0.8, 1.2) \text{ fm}$ for model a or $(R_S, R_L) = (0.7, 1.0) \text{ fm}$ for model b [36,37]. The different fitting procedures result in different values for the 26 unknown LECs governing the strength of short-range terms in the interactions. Accompanying these 2N interactions is the leading chiral 3N interaction which introduces two unknown LECs $c_D$ and $c_E$ (in standard notation) constrained to reproduce the trinucleon binding energies and, concurrently, either the GT matrix element contributing to tritium $\beta$-decay [39] in the starred model or the nd-doublet scattering length in the nonstarrned one [40].

Lastly, the vector and axial current operators entering the calculation were derived with time-ordered perturbation theory by the JLab-Pisa group using the same $\chi$EFT formulation as the NV2+3 interactions. Details about the electroweak currents used in this Letter can be found in Refs. [24,25,38,39,52–54].

**Results.** The results of the VMC calculation of the partial muon capture rate in $A = 3$ and $A = 6$ using the NV2+3 nuclear Hamiltonian are presented in Table I. Capture rates were determined using nuclear axial and vector current operators consistent with the NV2+3 model. The nuclear axial currents [38] contain only tree-level diagrams whereas the vector current operators account for loop corrections derived in Refs. [24,25,53,54].

Calculations of the rate with the leading order one-body only [$\Gamma(1b)$] and one- plus two-body electroweak currents [$\Gamma(2b)$] were performed for ground state to ground state transitions. The partial capture rate on $^3$He has been precisely measured [55], and the one-body contribution alone cannot reproduce this measurement. With the two-body electroweak currents included, the VMC rates increase by about 9%–16%. At this level, the agreement with the datum ranges from about 0.1%–4.6%. How the 3N interaction was fit has the most significant impact on the rate, leading to differences on average of 54 s$^{-1}$ whenever the 3N interaction is changed. Note that the LEC $c_D$ entering the 3N interaction governs the strength of the axial contact current at next-to-next-to-next-to-leading order in the chiral expansion [51]. Therefore, variations in the 3N interaction lead to variations in the current as also observed in the study of Ref. [51] on $\beta$-decay matrix elements. The cutoff and energy range of the fit lead to changes in the rate of 16 and 22 s$^{-1}$ on average, respectively, which is consistent with the findings of Refs. [12,15].

In the $^6$Li capture, the inclusion of two-body electroweak currents also increases the rate with a greater enhancement in the nonstarrned models relative to their starred counterparts. Even with this increase, ranging approximately from 3% to 7%, the rates predicted at the VMC level for the NV2+3 models are about 11–21% slower than the available experimental datum [56]. Here, the difference due to the 3N

| Capture | Model | $R$ | $\Gamma(1b)$ (s$^{-1}$) | $\Gamma(2b)$ (s$^{-1}$) | Expt. (s$^{-1}$) |
|---------|-------|----|----------------------|----------------------|----------------|
| $^3$He($^{12}$+, $^1\bar{\nu}$) $\rightarrow$ $^3$He($^{12}$+, $^1\bar{\nu}$) | Ia (lb) | 0.995 (0.995) | 1350.3 ± 0.8 (1363.4 ± 0.2) | 1564.4 ± 0.9 (1545.7 ± 0.3) | 1496.0 ± 4.0 [55] |
| | Ia* (Ib*) | 0.995 (0.995) | 1357.4 ± 0.2 (1358.5 ± 0.2) | 1473.9 ± 0.3 (1483.6 ± 0.3) | |
| | Iia (Iib) | 0.995 (0.995) | 1369.7 ± 0.2 (1372.8 ± 0.2) | 1533.8 ± 0.3 (1512.4 ± 0.3) | |
| | Iia* (Iib*) | 0.995 (0.995) | 1364.5 ± 0.2 (1372.5 ± 0.2) | 1484.4 ± 0.3 (1497.0 ± 0.3) | |
| $^6$Li($^{14}$+, 0) $\rightarrow$ $^6$He(0+, 1) | Ia (lb) | 0.990 (0.990) | 1196 ± 2 (1243 ± 2) | 1282 ± 2 (1331 ± 2) | 1600$^{+330}_{-129}$ [56] |
| | Ia* (Ib*) | 0.990 (0.990) | 1154 ± 3 (1188 ± 2) | 1177 ± 3 (1233 ± 2) | |
| | Iia (Iib) | 0.990 (0.990) | 1227 ± 2 (1142 ± 2) | 1294 ± 2 (1185 ± 2) | |
| | Iia* (Iib*) | 0.990 (0.990) | 1215 ± 2 (1151 ± 2) | 1257 ± 2 (1185 ± 2) | |
interaction is no longer the dominant contribution to the uncertainty. We find that, on average, the cutoff and energy range of the fit both change the rate by 72 s⁻¹, whereas the 3N interaction changes the rate by 60 s⁻¹.

We compute a VMC average for both rates under study and use the average changes due to the chiral 2N interaction cutoffs, the energy range used to fit the interaction, and the 3N interaction to assign a total error bar. An additional source of uncertainty was considered by varying the nucleonic axial radius parameter by ±10%. We found that, on average, the difference in the rate was ±0.6% due to this variation. We combine the four uncertainties in quadrature to determine the overall uncertainty on the VMC averages, obtaining Γ(A=3; VMC) = 1512 ± 32 and Γ(A=6; VMC) = 1243 ± 59 s⁻¹.

In addition to the VMC calculation, a GFMC propagation was performed for models Ia and Ia*, and corresponding results are reported in Table II. These two models provided the fastest and slowest VMC partial capture rates for A = 3 and should give an upper and a lower limit on GFMC rates. The GFMC method removes spurious contamination from the VMC wave functions by propagating them in imaginary time and should, thus, provide more reliable results for these two nuclear Hamiltonians. Figure 1 displays our average VMC results as well as both VMC and GFMC results for models Ia and Ia*, compared with experimental data and past theoretical calculations. The GFMC error is taken to be half of the difference between the two available calculations.

At the VMC level, model Ia overpredicted the A = 3 muon capture rate by 4.6%. After propagation, the rate is decreased and reaches a 1.5% agreement with the datum. By contrast, model Ia*, which had 1.4% agreement with the experimental datum at the VMC level, now underpredicts the rate by 4.2%.

In Fig. 1(a), one sees that the results of past chiral calculations in Refs. [12,15] fall within the bounds for the NV2+3 GFMC rate provided by models Ia and Ia*. Even when using the more recent value of Gv, the rate of Ref. [12] falls within our GFMC band. Although there is this agreement, because these past calculations use a different set of chiral currents and underlying nuclear interactions than the present Letter, it is difficult to directly compare them to our GFMC results. In the future, benchmark calculations with other ab initio methods based on the same dynamical inputs would be useful to further validate our microscopic approach.

Whereas the A = 3 GFMC rates exhibit few percent decreases from the VMC ones, the A = 6 rates display a dramatically different behavior for models Ia and Ia*. The matrix elements for the model Ia calculation were fairly stable when propagated from VMC to GFMC, resulting in a modest subpercent change in the overall rate. However, for model Ia*, the dominant matrix elements changed at the few percent level, but since the rate is proportional to the square of the matrix element, this leads to a roughly 20% change in the rate.

To further understand this behavior, one can look at the system size as a function of τ during the GFMC propagation of the A = 6 nuclei. The system size for ⁶Li(1⁺; 0) grows at the same rate in τ for both models; however, the ⁶He(0⁺; 1) ground state size is stable for model Ia, but increases monotonically in τ before beginning to converge for model Ia* (see the Supplemental Material [73]). Because of the e⁻iqτ dependence in the dominant one-body terms of the current operator, the matrix elements at a finite value of q experience a more significant drop for model Ia* due to the diffuseness of ⁶He(0⁺; 1) with that interaction. Performing the same analysis for the A = 3 system, we find that the system size is consistent between both models as a function of τ, explaining the similarity in their decreasing trend for this partial muon capture rate.

The difference with experiment in A = 6 is significant for both models Ia and Ia*, especially when compared with the few percent agreement obtained in GFMC calculations of the GT matrix element for the ⁶He → ⁶Li β-decay [51]. As detailed in Ref. [63], calculations of this rate [57–63] have ranged from 1160 to 1790 s⁻¹. The calculation of Ref. [63] matched the experimental datum by modeling ⁶Li as a ³He + τ

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**TABLE II.** VMC and GFMC calculations of partial muon capture rates in ³He and ⁶Li obtained with chiral one- and two-body axial and vector currents with the NV2+3 models. The experimental result is given in the last column. All uncertainties on the theoretical predictions are Monte Carlo errors.

| Capture | Model | Γ (VMC) (s⁻¹) | Γ (GFMC) (s⁻¹) | Expt. |
|---------|-------|---------------|----------------|-------|
| ³He(1⁺; ½) → ³H(0⁺; ½) | Ia    | 1564.4 ± 0.9  | 1519 ± 3       | 1496.0 ± 4.0 [55] |
|              | Ia*   | 1473.9 ± 0.3  | 1433 ± 2       |       |
| ⁶Li(1⁺; 0) → ⁶He(0⁺; 1) | Ia    | 1282 ± 2      | 1277 ± 10      | 1600 ± 330 [56] |
|              | Ia*   | 1177 ± 2      | 926 ± 8        |       |

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**Fig. 1.** The partial muon capture rate in (a) ³He and (b) ⁶Li from the the NV2+3-Ia and NV2+3-Ia* models in VMC (light blue circle) and GFMC (dark blue star) calculations compared with other work (orange squares) [12–16,57–63]. The experimental values (dashed gray line) and their errors (shaded region) [55,56] are included for comparison with the theory predictions.
cluster and using the Fujii-Primakoff [57] effective Hamiltonian for muon capture. Subpercent agreement was also obtained by Ref. [59], which treated the $^6$Li and $^6$He nuclei as elementary particles with magnetic and axial form factors extracted from experiment. The two calculations presented by the authors of that work adopted different formulations of the partially conserved axial current relation to obtain the pseudoscalar form factor with the faster rate using the Gell-Mann-Lévy version [74] and the slower rate using the Nambu one [75]. The Nambu definition is consistent with the induced pseudoscalar term in the weak axial current from $\chi$EFT. It is difficult to compare our result with those of other theoretical treatments of the $^6$Li partial capture rate, particularly, since most of these treatments are decades old. For example, in the work of Ref. [61] the weak-interaction Hamiltonian is that of Eq. (1); however, the $^6$Li and $^6$He bound states are described by shell model wave functions with valence configurations restricted to the 1p shell; moreover, the nuclear electroweak current neglects meson-exchange contributions [61]. We find that our result at leading order (obtained with one-body currents) is quenched relative to the shell model one as we would have expected (see Ref. [51]). More modern calculations with other $ab\ initial\ initio$ methods and a novel measurement of the rate would be valuable in establishing the validity of our nuclear inputs and the many-body approach.

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