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Physics of eccentric binary black hole mergers: A numerical relativity perspective

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Gravitational wave observations of eccentric binary black hole mergers will provide unequivocal evidence for the formation of these systems through dynamical assembly in dense stellar environments. The study of these astrophysically motivated sources is timely in view of electromagnetic observations, consistent with the existence of stellar mass black holes in the globular cluster M22 and in the Galactic center, and the proven detection capabilities of ground-based gravitational wave detectors. In order to get insights into the physics of these objects in the dynamical, strong-field gravity regime, we present a catalog of 89 numerical relativity waveforms that describe binary systems of non-spinning black holes with mass-ratios $1 \leq q \leq 10$, and initial eccentricities as high as $e_0 = 0.18$ fifteen cycles before merger. We use this catalog to quantify the loss of energy and angular momentum through gravitational radiation, and the astrophysical properties of the black hole remnant, including its final mass and spin, and recoil velocity. We discuss the implications of these results for gravitational wave source modeling, and the design of algorithms to search for and identify eccentric binary black hole mergers in realistic detection scenarios.

I. INTRODUCTION

The gravitational wave (GW) detection of several binary black hole (BBH) mergers [1–6], and the first multi-messenger observation of two colliding neutron stars (NSs) in gravitational and electromagnetic waves [7] have shed light into the nature of gravity in the most extreme astrophysical settings, and has unveiled the identity of the central engines that power the most energetic electromagnetic explosions in the Universe [8–11], while also providing the means to put at work visionary methods to use GWs to quantify the rate of expansion of the Universe [12–15].

Along this trail of discovery, it has also become evident that numerical relativity (NR) plays a central role to understand the physics of GW sources, and to inform the development of signal-processing algorithms to detect and characterize these astrophysical events [16–21], and astrophysical sources that still await discovery [22–30].

In preparation for the characterization of BBH mergers whose astrophysical properties span a parameter space that has not yet been probed by existing GW detections, several NR groups are working in earnest to construct large-scale NR waveform catalogs [17, 37, 38]. Since these activities have thus far focused on the study of quasi-circular BBH mergers, in this article we fill in a critical void in the literature by presenting a comprehensive study of the physics of moderately eccentric BBH mergers.

The rationale for this study is multifold. From the perspective of electromagnetic observations, recent findings are consistent with the existence of stellar-mass BHs in the vicinity of the Galactic center, and in the Galactic Cluster M22 [39–42]. These observations have triggered the development of numerical models that provide a realistic description of the formation and retention of BBHs in dense stellar environments, correcting previous calculations based on $N$-body simulations that did not include post-Newtonian corrections [43] to model the orbital dynamics of these systems, thereby underestimating the merger rates of these systems by orders of magnitude [42, 44–79]. In summary, we have evidence for the existence of stellar-mass BHs that may form eccentric compact binary systems in dense stellar environments, and consequently be detected through GW emission. Through this study, we provide new insights into the physics of these astrophysically motivated sources.

Furthermore, as discussed in [80], no matched-filtering algorithm has been presented in the literature that is tailored for the detection of $\ell = |m| = 2$ eccentric waveforms [80]. However, signal processing algorithms based on deep neural networks have been used to demonstrate that moderately eccentric BBH mergers can be detected and characterized from real LIGO noise, considering both NR waveforms that only include the leading order quadrupole mode $\ell = |m| = 2$ [22–25], and higher-order waveform multipoles [27]. We expect that this NR waveform catalog may be used to quantify the sensitivity of burst searches, and of next-generation neural network models that are tailored to detect and characterize eccentric BBH mergers. In summary, activities around modeling, detection and characterization of eccentric BBHs
are reaching the required level of maturity to establish or rule out the existence of compact binary populations in dense stellar environments.

To advance our understanding of the physics of compact binary populations in dense stellar environments, in this article we introduce a NR waveform catalog that describes eccentric BBH mergers, and utilize it to get insights into the dynamics of these GW sources, e.g., the energy and angular momentum loss through GW emission, and the astrophysical properties of the BH remnant, i.e., its final mass and spin as a function of initial eccentricity and mass-ratio, as well as the recoil velocity of the BH remnant. These studies will inform ongoing GW modeling efforts, and the development of signal-processing algorithms to search for and identify these sources. This article is organized as follows. Section II describes the properties of our NR catalog. In Section III we compute the energy and angular momentum radiated away through GW emission, making pair-wise comparisons between NR waveforms that include the \( \ell = |m| = 2 \) mode or higher-order waveform multipoles. In Section IV we compute the astrophysical properties of the BH remnants, and compare these results with those obtained for quasi-circular BBH mergers. We describe the relevance of these analyses in terms of GW modeling efforts for eccentric BBH mergers in Section V. We summarize our findings and outline future directions of work in Section VI.

II. NUMERICAL RELATIVITY CATALOG

We have produced a catalog of 89 simulations with the open source, NR software, the Einstein Toolkit [81–100]. This catalog describes non-spinning BBHs with mass-ratios \( 1 \leq q \leq 10 \) and eccentricities as high as \( e_0 = 0.18 \) fifteen cycles before merger. A visualization of this catalog may be found at [101, 102]. We have post-processed the data products of these simulations using the open source software stack POWER [103], and extracted the modes \((\ell, |m|) = \{(2, 2), (2, 1), (3, 3), (3, 2), (3, 1), (4, 4), (4, 3), (4, 2), (4, 1)\}\). As described in Appendix B, each of these simulations was produced with several levels of resolution to quantify convergence. The real part of the \( \ell = |m| = 2 \) mode, extracted at future null infinity, for each NR waveform is presented in Figure 1. The properties of these NR waveforms are listed in Table II.

Characterizing the properties of the NR waveforms presented in Table II requires the construction of a method to quantify the orbital eccentricity of these simulations. Using the orbital evolution of these simulations to obtain an estimate of the orbital eccentricity is inadequate due to the gauge-dependent nature of the binary’s orbit. On the other hand, methods to construct initial data for spinning BHs on quasi-circular orbits have also introduced definitions of orbital eccentricity, based on orbital separations and waveform phase and amplitude of the Weyl scalar \( \psi_4 \) [104]. However, while the scope of the method introduced in [104] is to construct high-quality initial data for quasi-circular mergers, and therefore, using \( \mathcal{O}(e) \) approximations to model the effect of eccentricity may suffice, we aim to measure larger values of orbital eccentricity.

To address this matter, we have used the inspiral-merger-ringdown ENIGMA waveform model introduced in [26] to determine the eccentricity, mean anomaly and gauge-invariant frequency parameters, \((e_0, \ell_0, x_0)\), that optimally describe each NR waveform in our catalog. We do this by finding the \((e_0, \ell_0, x_0)\) triplet that maximizes the overlap between each NR waveform and its ENIGMA counterpart. In [105] we quantified the optimal time window to remove junk radiation while keeping intact the signatures of eccentricity at early times in the NR waveforms. Such time range is given by \( t \leq 60M \). A detailed description of this method, including the corresponding open source software stack for its use to characterize NR waveforms catalogs at scale, is presented in an accompanying article [105].

In brief, we construct our method using the inspiral evolution of the ENIGMA waveform model, which contains state-of-the-art post-Newtonian corrections for eccentric binaries, which include eccentricity corrections in the conservative and radiative pieces up \( \mathcal{O}(e^{12}) \), including instantaneous, tails and tails-of-tails contributions, and a contribution due to nonlinear memory; and quasi-circular corrections both from post-Newtonian, self-force and perturbative calculations up to \( \mathcal{O}(x^6) \) [26, 29]. Furthermore, in [26], we have demonstrated that \( e_0 = 0 \) ENIGMA waveforms capture the dynamics of quasi-circular BBH mergers with excellent accuracy. We showed this by computing overlaps between quasi-circular ENIGMA waveforms and their quasi-circular Effective One Body (EOB) counterparts [106]. Assuming advanced LIGO’s Zero Detuned High Power sensitivity [107], and using an initial GW frequency of 15 Hz to compute the overlaps, Figure 2 in [26] shows that the overlap between quasi-circular ENIGMA and EOB waveforms is \( \mathcal{O} \geq 0.99 \). Since the waveforms we are characterizing in this study are much shorter than those used to assess the accuracy of the ENIGMA model in the quasi-circular limit, it follows that ENIGMA will capture the dynamics of moderately eccentric systems with excellent accuracy.

It is worth highlighting that while the ENIGMA waveform model was originally validated with eccentric NR waveforms that describe BBH mergers with mass-ratios \( q \leq 5.5 \) and eccentricities \( e_0 \leq 0.18 \) twenty cycles before merger [26], it is through this analysis, and with the availability of new NR waveforms, that we can now report that the ENIGMA model can accurately describe BBH mergers with mass-ratios up to \( q = 10 \) with \( e_0 \leq 0.18 \) fifteen cycles before merger.
FIG. 1. For a given mass-ratio \( q \), each row presents the real part of the \( \ell = |m| = 2 \) mode of each waveform in our catalog, extracted at future null infinity. The initial eccentricity, \( e_0 \), increases from left to right. All these waveforms have unit amplitude at \( t = 0 \). Table II lists the properties of this waveform catalog.
III. ENERGY AND ANGULAR MOMENTUM EMISSION OF ECCENTRIC BLACK HOLE MERGERS

For each NR waveform in our catalog, we have quantified the energy, \( E \), and angular momentum, \( J \), radiated away through GW emission using the relations [108]

\[
\Delta E = \frac{1}{16\pi} \int_{t_0}^{t_f} \sum_{\ell = 0}^{\ell_{\text{max}}} \sum_{m = -\ell}^{\ell} dt' |N_{\ell m}(t')|^2 , \tag{1}
\]

\[
\Delta J = \frac{1}{16\pi} \int_{t_0}^{t_f} \sum_{\ell = 0}^{\ell_{\text{max}}} \sum_{m = -\ell}^{\ell} dt' m \Im [h_{\ell m}(t') N_{\ell m}^*(t')] , \tag{2}
\]

\[
N_{\ell m}(t) = \frac{d h_{\ell m}(t)}{dt} , \tag{3}
\]

where \( N_{\ell m}(t) \) represents the complex news function at infinity. The integration is done from the time the NR waveform is free from junk radiation, \( t_0 = 60M \), to the final sample time of the NR waveform, \( t_f \). For these calculations we have considered the \((\ell, |m|) = \{(2, 2), (2, 1), (3, 3), (3, 2), (3, 1), (4, 4), (4, 3), (4, 2), (4, 1)\} \) modes. It is worth pointing out that the choice \( t_0 \to 60M \) is informed by the study presented in [108], which demonstrated that this choice removes high-frequency noise while keeping intact the imprints of eccentricity in the NR waveforms once they are free from junk radiation. Using Eqs. (1)-(3), in Figure 2 we quantify the importance of including higher-order waveform modes to compute the energy and angular momentum carried away by GWs. We do this through pair-wise comparisons between NR waveforms that include either all the modes listed above, or just the \( \ell = |m| = 2 \) mode, using the relations

\[
\Delta E' = \frac{\Delta E(\ell, |m|) - \Delta E(\ell = |m| = 2)}{\Delta E(\ell, |m|)} , \tag{4}
\]

\[
\Delta J' = \frac{\Delta J(\ell, |m|) - \Delta J(\ell = |m| = 2)}{\Delta J(\ell, |m|)} . \tag{5}
\]

Using the two highest resolution runs for each simulation in our catalog, we computed \((\Delta E', \Delta J')\), and found that the largest difference between these two independent measurements is \( \leq 5\% \). The values we present in Figure 2 were extracted from the highest resolution runs. We notice that for each mass-ratio BBH population, i.e., if we consider a given set of markers in the panels of Figure 2, \((\Delta E', \Delta J')\) are nearly constant across the eccentricity range that we have explored in this study. In different words, \((\Delta E', \Delta J')\) are constant polynomials in eccentricity for \( e_0 \leq 0.2 \). We observe a minor deviation from this pattern at the high-end of the eccentricity range for the most asymmetric mass-ratio BBH systems. To be precise, if we fit a constant polynomial using the two lowest eccentricity samples for each mass-ratio population, we find that the largest deviation occurs for the most eccentric sample of the \( q = 10 \) BBHs, with a fractional error \( \leq 8\% \) for the measurement of \( \Delta J' \).

These results also show that that for systems with \( q \geq 5 \) it is essential to include higher-order waveform modes to accurately describe the dynamics of eccentric BBH mergers. This result is consistent with recent studies [27], which indicate that the inclusion of higher-order modes for non-spinning, eccentric BBH mergers has a more significant impact for GW detection, in the context of signal-to-noise ratio calculations, than for their non-spinning, quasi-circular BBH counterparts. It is worth highlighting that the eccentric NR waveform we have produced for this analysis for \( q \geq 5 \) are the first of their kind.
in the literature, so these results shed new light on the importance of including higher-order waveforms modes for the modeling of radiated energy and angular momentum of eccentric BBH mergers.

IV. FINAL MASS, SPIN, AND RECOIL VELOCITY OF POST-MERGER BLACK HOLES

We have computed the final mass, \( M_f \), and final spin, \( q_f \), of the BH remnant using the QuasiLocalMeasures thorn of the Einstein Toolkit. The final mass is given by

\[
M_f = \sqrt{M_{irr}^2 + \frac{q_f^2}{4M_{irr}}}, \quad \text{where (6)}
\]

\[
M_{irr} = \sqrt{\frac{A}{16\pi}}, \quad \text{(7)}
\]

\( M_{irr} \) is the irreducible mass, given in terms of the BHs’ event horizon area, \( A \). \( q_f \) is computed as the Komar angular momentum \([109–112]\)

\[
q_f = \frac{1}{8\pi} \oint_S K_{ij} s^i \phi^j dA,
\]

where the integral is over the surface, \( S \), of the apparent horizon, \( K_{ij} \) is the extrinsic curvature, \( s^i \) is a spacelike, outward normal to the horizon, and \( \phi^i \) is a Killing vector associated with the rotational symmetry around the spin axis.

In Figure 3 we present results for \( (M_f, q_f) \). As before, we used our two highest resolution runs for each NR simulation to compute these observables, and found that these two independent measurements differ by \( \leq 3\% \). The results presented for \( (M_f, q_f) \) in Figure 3 were extracted from the highest resolution runs in our catalog. We notice that for each mass-ratio population, i.e., for a given set of markers, the final mass and spin of the BH remnant are nearly independent of eccentricity in the range \( e_0 \leq 0.2 \). We conclude this since both \( (M_f, q_f) \) can be described as constant polynomials in eccentricity within the range we have considered in this study. We can directly compare these results using formulae derived for quasi-circular BBH mergers in \([113, 114]\). Notice that we have included horizontal gray lines in both panels that provide the predictions for \( (M_f, q_f) \) in the \( e_0 \to 0 \) limit.

For the \( M_f \) results (top panel in Figure 3) the gray lines present the quasi-circular prediction for the final mass of the BH remnant for the mass-ratios \( q = \{1, 1.5, 2, 2.5, 3, 3.5, 4, 4.5, 5, 6, 8, 10\} \) from bottom to top, respectively. We notice that the equal mass-eccentric BBH population presents the largest deviation from the quasi-circular prediction. However, this discrepancy is \( \leq 1\% \). In the case of the final spin of the BH remnant, the bottom panel of Figure 3 also presents the quasi-circular predictions for this observable. Notice, however, that in this case, the grey lines describe the mass-ratios listed above but now from top to bottom. As in the case of \( M_f \), our results for the final spin of moderately eccentric BBH mergers are fairly consistent with results obtained from quasi-circular BBH mergers. This can only be the case if the eccentric NR waveforms we have produced in this catalog circularize prior to merger. We have explored this scenario in detail, and have found that this is indeed the case. For a sample case, Figure 4 presents two waveform signals produced by BHs that have the same separation, but different initial eccentricity. The eccentric waveform contains all the telltale signatures of eccentricity, i.e., significant modulations in the amplitude and phase at early times, which correspond to periape (local maxima) and apoapse (local minima) passages. We also observe that the waveform circularizes very rapidly, from \( e_0 = 0.18 \) fifteen cycles before merger, turning into a quasi-circular waveform signal near the merger event. This is the reason why the results presented in Figure 3 are consistent with their quasi-circular counterparts.

Earlier work on this front includes \([115]\), which presented calculations for the final spin and circularization of equal-mass eccentric BBH mergers, and showed that for BBH mergers with larger initial eccentricities than those considered in this work, the final spin of the BH remnant is greater than its quasi-circular counterpart. Additionally, reference \([31]\) discussed the circularization of moderately eccentric BBH mergers with \( q \leq 3 \). In this article, we provide a systematic study of the observables \( (q_f, M_f) \) to furnish evidence for the circularization of eccentric BBH mergers with \( q \leq 10 \) and \( e_0 \leq 0.18 \) fifteen cycles before merger.

We have also computed the recoil velocity of eccentric BBH mergers, \( |v|_{\text{kick}} \), using the PunctureTracker thorn in the Einstein Toolkit. To do this, we have considered the last 100M of evolution of our NR simulations, and a simple first order finite difference formula for the velocity in terms on the measured locations. In Table I we present the minima and maxima of the recoil velocity, \( |v|_{\text{kick}} \), for the range of eccentricities we consider for each mass-ratio. We have also obtained gauge-invariant perturbations following \([118]\) (see Eqs. (33)–(39) therein). These gauge-invariant results are presented in the last column of Table I.

Important observations to be drawn from Table I include: (i) the kick velocity of BH remnants produced by quasi-circular BBH mergers, \( |v|_{\text{kick}}^{q=0} \), was obtained using the formulae presented in \([116]\); (ii) the recoil velocity for all the \( q = 1 \) eccentric BBH mergers in our catalog is \( |v|_{\text{kick}} = 0 \), which is consistent with results obtained for non-spinning, quasi-circular BBH mergers \([117]\); (iii) the kick velocities for our \( 1 \leq q \leq 10 \) population of eccentric BBH mergers are fairly consistent with the expected values of their quasi-circular counterparts, even though the formulae used to estimate \( |v|_{\text{kick}}^{q=0} \) was calibrated with quasi-circular BBH mergers with mass-ratios \( q \leq 8 \).
These results for \((M_f, q_f, |v|_{\text{kick}})\) cover an entirely new region of parameter space in the modeling of eccentric BBH mergers, providing new insights into the physics of these GW sources. We discuss the implications of these findings in the following section.

\section{Implications for the Modeling and Detection of Eccentric Mergers}

To date, there are only a handful of inspiral-merger-ringdown waveform models that describe the GW emission of eccentric BBH mergers \cite{2000, 2002, 2003, 2004, 2005}. These models assume that moderately eccentric BBHs circularize prior to the merger event. This assumption is sound, in light of the results presented in the previous section, for BBHs with \(q \leq 10\) and whose residual eccentricity is as high as \(e_0 \leq 0.18\) just fifteen cycles before merger.

Furthermore, we have found that for the most extreme sample of our NR catalog, e.g., P0204, which represents BBHs with \(q = 10\) and \(e_0 = 0.18\) fifteen cycles before merger, circularization is only attained right before merger, as shown in Figure 4. In different words, while assuming circularization of moderately eccentric BBH mergers is a reasonable \textit{ansatz}, this also means that the modeling of these GW sources demands the development of an inspiral evolution scheme that provides an accurate description of the dynamical evolution of these objects throughout the inspiral evolution, and which remains accurate one or two cycles before merger. To accomplish this level of accuracy so late in the inspiral evolution,
VI. CONCLUSION

We have studied the physics of eccentric BBH mergers using a NR waveform catalog that describes BBH systems with mass-ratios $q \leq 10$, and initial eccentricities $e \leq 0.18$ up to fifteen cycles before merger.

We quantified the importance of including higher-order waveform modes to compute the energy and angular momentum carried away by GWs in eccentric BBH mergers. We have also demonstrated that the properties of BH remnants described by our NR catalog are consistent with their quasi-circular counterparts, which provides evidence for the circularization of moderately eccentric BBH mergers. We have also computed recoil velocities of BH remnants produced by eccentric BBH mergers, and found that these are fairly similar to those computed in the literature for non-spinning, quasi-circular BBH mergers.

Based on these analyses, we have provided evidence that existing source modeling efforts that assume the circularization of moderately eccentric BBH mergers is sound. We have also shown that since circularization takes place close to the merger event, any semi-analytical model that is used to describe these GW sources should include higher-order corrections both to the conservative and radiative pieces of the source’s dynamics, and to the waveform strain.

Recent studies in the literature have identified parameter space degeneracies between orbital eccentricity and spin corrections [29]. In order to get better insights into this finding, it is essential to understand the dynamics of spinning BHs on eccentric orbits, and then use the NR waveform catalog we have introduced in this study to carefully assess in which regions of parameter space such a degeneracy may be broken to distinguish these two compact binary populations.

The construction of a NR waveform catalog for spinning BHs on eccentric orbits is already underway to shed light on this timely, and astrophysically motivated study. Specific aspects to address in such study will encompass: (i) orbital configurations that significantly shorten the length of waveform signals, i.e., eccentricity and spin anti-aligned configurations; (ii) competing effects to determine the length of waveform signals, i.e., rapidly spinning BHs on spin-aligned configurations (which increase the length of waveforms as compared to non-spinning BBHs) vs moderate values of initial eccentricity (which decrease the length of waveforms as compared to quasi-circular BBHs); and (iii) identify telltale signatures of GW sources that can be used to infer the existence of eccentric compact binary populations through GW observations, e.g., astrophysical properties of the BH remnant, and the coupling of eccentricity and spin-spin and spin-orbit effects at peri-apse passages during the inspiral evolution of these systems.

The modes we have extracted in this study with the open source POWER package do not include $m = 0$ memory modes. Their extraction requires the use of the Cauchy Characteristic Extraction method [148]. Given the importance of these modes for the characterization of eccentric BBH mergers, we will present these modes in a forthcoming study, accompanied by a systematic analysis on the observability of these modes with second and third generation GW detectors.

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Appendix A: Properties of Numerical Relativity Catalog

Table II lists the properties of our numerical relativity catalog.

| Simulation | q | $e_0$ | $\ell_0$ | $x_0$ |
|------------|---|-----|-----|-----|
| E0001      | 1 | 0.052 | 3.0  | 0.070 |
| E0009      | 2 | 0.052 | 3.0  | 0.0794|
| E0013      | 2.5| 0.050 | 3.0  | 0.0813|
| E0017      | 3 | 0.050 | 3.0  | 0.0831|
| F0002      | 3 | 0.066 | 3.0  | 0.0780|
| F0010      | 2 | 0.066 | 3.0  | 0.0803|
| F0014      | 2.5| 0.068 | 3.0  | 0.0822|
| F0018      | 3 | 0.068 | 3.0  | 0.0842|
| G0003      | 1 | 0.094 | 3.0  | 0.0788|
Appendix B: Convergence of the numerical waveforms

We use a grid setup based on the setup used in [89]. There is a central, mesh refined cubical region of the grid in which Cartesian coordinates are used, surround by 6 regions that make up a cubed sphere grid with constant angular resolution.

We use 8th order finite differencing operators to compute spatial derivatives of the spacetime quantities in the Einstein field equations. This requires the use of 5 ghost zones, and together with using a classical 4th order Runge-Kutta timestepper implies that each refined region is surrounded by 20 points that are filled in via prolongation from the next coarser region. We use vertex centered 5th order prolongation operators rather than full 8th order prolongation operators.

The cubical region employs mesh refinement with the resolution on the coarsest grid being $h_{\text{coarse}} = 1.92 \, M$. Each of the black holes is surrounded by a set of nested moving boxes such that the resolution in the finest box containing the black hole $i$ is $1.2 M_i/(N_i - 1)$ where $M_i$ is the initial mass parameter of black hole $i$ and $N_i$ is the number of points used for the resolution level $l$ simulation. In our simulations we used $N_i = 32, 36, 40, 44$, where $N_i = 44$ was only used for simulations with a mass ratio $q > 5$. The finest box surrounding each black hole has a radius of $1.2 M_i$ and each coarser box has twice the radius of the next finer one. During the simulation we track the location of each black hole and keep the set of nested refined boxes approximately centered on the black hole. Finally the outer edge of the cubical region is chosen large enough to contain all refined regions including their prolongation regions.

In the spherical region we choose an angular resolution of $h_{\text{angular}} = \pi/(4 N_i)$ and a radial resolution of $1.92 \, M$ which matches the coarsest resolution in the Cartesian grid. The outer boundary is chosen such that it is causally disconnected from the outermost detector at which we extract gravitational waves from.

We use a time step $\Delta t = 0.864 \, M$ on the coarsest level, corresponding to a Courant-Friedrichs-Lewy condition of $\Delta t/h_{\text{coarse}} = 0.45$ which is held constant on the finer levels by decreasing their time step size.

We extract gravitational waves using modes of the Weyl scalar $\psi_4$ extracted on coordinate spheres of radius $r_{\text{det},i} = 100 \, M, 115 \, M, 136 \, M, 167 \, M, 214 \, M, 300 \, M, 500 \, M$.

Using 8th order finite differencing operators our simulations would, under ideal circumstances, converge towards the correct solution with an error term which scales like $h^8$, where $h$ is the spatial resolution of the simulation. However due to lower order schemes present in the simulation, for example the interpolation at mesh refinement boundaries which is only 5th order accurate, as well as artifacts caused by the adaptive mesh refinement logic making independent decisions where to refine for each simulation, the observed convergence order typically differs from 8.

To estimate the convergence of each waveform we simulated each set of physical parameters using at least 3 (4) simulations using increasing resolution for waveforms of mass ratio $q \leq 5$ ($q > 5$). We then compute the gravitational wave phase $\phi^{(2,2)}$ from the complex $\ell = m = 2$ mode of the spherical harmonic decomposition of the outgoing component of the Weyl scalar $\psi_4$ and studied its convergence properties.

Figure 5 shows the rescaled phase differences $\phi^{(2,2)}(h_n) - \phi^{(2,2)}(h_{\text{high}})$ between the gravitational wave phase obtained from the simulation with resolution $h_n$ and the highest resolved simulation. Phase differences have been rescaled such that for a convergent simulation the plotted curves overlap. For case J0046 we observe an approximate convergence order of $N \approx 4.5$ which is within the range of expected values.

Not all simulated cases show clean convergence behaviour, with the convergence order for some of them being larger than 8, which may indicate that our lowest resolution simulation does not adequately resolve the features present in the simulation domain, and others swapping the ordering of phases between the low, medium and high resolution simulations, making an estimate of the convergence order impossible.

Given that large number of simulations, this is to be

| L0014 | 5.0  | 0.080 | 3.0 | 0.0856 |
| L0015 | 5.0  | 0.100 | 3.0 | 0.0853 |
| L0016 | 5.0  | 0.140 | 2.9 | 0.0862 |
| L0017 | 5.0  | 0.060 | 3.0 | 0.0869 |
| L0018 | 5.0  | 0.068 | 3.0 | 0.0878 |
| L0019 | 5.0  | 0.100 | 3.0 | 0.0869 |
| L0020 | 5.0  | 0.140 | 2.9 | 0.0882 |
| L0029 | 4.5  | 0.058 | 3.0 | 0.0844 |
| L0030 | 4.5  | 0.080 | 3.1 | 0.0835 |
| L0031 | 4.5  | 0.120 | 3.1 | 0.0827 |
| L0032 | 4.5  | 0.180 | 3.0 | 0.0849 |
| L0033 | 5.0  | 0.060 | 3.0 | 0.0852 |
| L0034 | 5.0  | 0.080 | 3.0 | 0.0852 |
| L0037 | 5.0  | 0.060 | 3.0 | 0.0870 |
| L0038 | 5.0  | 0.080 | 3.0 | 0.0870 |
| L0039 | 5.0  | 0.120 | 2.9 | 0.0867 |
| L0040 | 5.0  | 0.180 | 2.9 | 0.0864 |
| P0001 | 6.0  | 0.050 | 3.0 | 0.0867 |
| P0004 | 6.0  | 0.140 | 2.9 | 0.0867 |
| P0006 | 8.0  | 0.080 | 2.9 | 0.0931 |
| P0007 | 8.0  | 0.100 | 2.9 | 0.0926 |
| P0008 | 8.0  | 0.140 | 2.9 | 0.0910 |
| P0009 | 10.0 | 0.060 | 2.9 | 0.0971 |
| P0013 | 6.0  | 0.054 | 3.0 | 0.0871 |
| P0014 | 6.0  | 0.078 | 2.9 | 0.0885 |
| P0016 | 6.0  | 0.160 | 2.8 | 0.0900 |
| P0017 | 8.0  | 0.060 | 3.0 | 0.0927 |
| P0020 | 8.0  | 0.180 | 2.9 | 0.0936 |
| P0022 | 10.0 | 0.080 | 2.9 | 0.0979 |
| P0023 | 10.0 | 0.120 | 2.9 | 0.0968 |
| P0024 | 10.0 | 0.180 | 3.0 | 0.0957 |
FIG. 5. Convergence of the phase difference in the gravitational wave phase for case J0046, rescaled to demonstrate convergence at order $N \approx 4.5$. We compute convergence in the time interval $500 \, M \leq t \leq t_{\text{max}} - 200 \, M$, where $t_{\text{max}}$ is the point of maximum amplitude, approximately corresponding to the time of merger. Note that this plot includes the pulse of junk radiation due to our initial data not containing any waves as well as the ringdown and merger signal, both of which are not convergent and thus lead to very large phase differences which are clipped in the plot.