Lamb Shift in Muonic Hydrogen

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Abstract
The Lamb shift in muonic hydrogen continues to be a subject of experimental and theoretical investigation. Here my older work on the subject is updated to provide a complementary calculation of the energies of the 2p-2s transitions in muonic hydrogen.

Introduction
The energy levels of muonic atoms are very sensitive to effects of quantum electrodynamics (QED), nuclear structure, and recoil, since the muon is about 206 times heavier than the electron [1]. In view of a proposed measurement of the Lamb shift in muonic hydrogen [2], an improved theoretical analysis seems to be desirable. Since the first theoretical analysis [3], the subject of the Lamb shift (the 2p-2s transition) in light muonic atoms has been investigated with increasing precision by a number of authors [4, 5, 6, 7, 8, 9, 10]. The present paper provides an independent recalculation of some of the most important effects, including hyperfine structure, and a new calculation of some terms that were omitted in the most recent literature, such as the virtual Delbrück effect [11]. An alternative calculation of the relativistic recoil correction is presented.

In the numerical calculations the fundamental constants from the CODATA 1998 ([12]) are used, i.e.: \( \alpha^{-1} \), \( \hbar c \), \( m_\mu \), \( m_e \), \( m_u = 137.0359998 \), 197.32696 MeV·fm, 105.658357 MeV, 0.5109989 MeV, 931.4940 MeV, respectively. The changes in these constants in the CODATA 2002 compared with CODATA 1998 are too small to make any relevant difference in the results.

Vacuum Polarization
The most important QED effect for muonic atoms is the virtual production and annihilation of a single \( e^+e^- \) pair. It has as a consequence an effective interaction of order \( \alpha Z \alpha \) which is usually called the Uehling potential ([13] [14]). This interaction describes the most important modification of Coulomb’s law. Numerically it is so important that it should not be treated using perturbation theory; instead the Uehling potential should be added to the nuclear electrostatic potential before solving the Dirac equation. However, a perturbative treatment is also useful in the case of very light atoms, such as hydrogen.
Unlike some other authors, we prefer to use relativistic (Dirac) wave functions to describe the muonic orbit. This is more exact, and as will be seen below, it makes a difference for at least the most important contributions. The wave functions are given in the book of Akhiezer and Berestetskii [15] and will not be given here. In perturbation theory, the energy shift due to an effective potential $\Delta V$ is given by

$$\Delta E_{n\kappa} = \frac{1}{2\pi^2} \int_0^\infty q^2 dq \Delta V(q) \cdot \int_0^\infty dr j_0(qr)[F^2_{n\kappa} + G^2_{n\kappa}]$$

where $F_{n\kappa}$ and $G_{n\kappa}$ are the small and large components of the wave function, $n$ is the principle quantum number and $\kappa$ is equal to $-(\ell + 1)$ if $j = \ell + \frac{1}{2}$ and $+\ell$ if $j = \ell - \frac{1}{2}$. $\Delta V(q)$ is the Fourier transform of the physical potential.

$$\Delta V(q) = 4\pi \cdot \int_0^\infty r^2 \cdot j_0(qr) \cdot \Delta V(r) dr$$

$$\Delta V(r) = \frac{1}{2\pi^2} \cdot \int_0^\infty q^2 \cdot j_0(qr) \cdot \Delta V(q) dq$$

As is well-known [1], the Uehling potential in momentum space is given by

$$V_{Ueh}(q) = -\frac{4\alpha(\alpha Z)}{3} \cdot G_E(q) \cdot F(\phi) = -4\pi(\alpha Z) \cdot G_E(q) \cdot U_2(q)$$

where $G_E$ is the proton charge form factor, $\sinh(\phi) = q/(2m_e)$ and

$$F(\phi) = \frac{1}{3} + (\coth^2(\phi) - 3) \cdot [1 + \phi \cdot \coth(\phi)]$$

$U_2(q)$ is defined in [1]. The vacuum polarization corrections were calculated in momentum space; the formulas (124,125,127) of [1] are completely equivalent to (200) in [10]. If the correction to the transition $2p_{1/2} - 2s_{1/2}$ is calculated in lowest order perturbation theory using nonrelativistic point Coulomb wave functions, the result is 205.0074 meV, in agreement with other authors [10].

The same procedure was used to calculate the two-loop corrections; the corresponding diagrams were first calculated by Källen and Sabry [16]. The Fourier transform of the corresponding potential is given in [1, 4]. The result for a point nucleus is 1.5080 meV.

In momentum space including the effect of nuclear size on the Uehling potential is trivial, since the corresponding expression for $\Delta V(q)$ is simply multiplied by the form factor. The numbers obtained were the same for a dipole form factor and for a Gaussian form factor, provided the parameters were adjusted to reproduce the experimental rms radius of the proton. The correction can be regarded as taking into account the effect of finite nuclear size on the virtual electron-positron pair in the loop. The contribution of the Uehling potential to the $2p-2s$ transition is reduced by 0.0081 meV with a proton radius of 0.862 fm [17], and by 0.0085 meV with a proton radius of 0.880 fm [18]. This result is consistent with the number given in [10] (eq.(266)). More recent values for the proton radius have been given by Sick [19] (0.895 ± 0.018 fm) and in the newest CODATA compilation [20] (0.875 ± 0.007 fm).
The numerical values given below were calculated as the expectation value of the Uehling potential using point-Coulomb Dirac wave functions with reduced mass:

|       | point nucleus | $R_p=0.875\text{fm}$ |
|-------|---------------|-----------------------|
|       | $2p_{1/2} - 2s_{1/2}$ | $2p_{3/2} - 2s_{1/2}$ | $2p_{1/2} - 2s_{1/2}$ | $2p_{3/2} - 2s_{1/2}$ |
| Uehling | 205.0282 | 205.0332 | 205.0199 | 205.0250 |
| Kaellen-Sabry | 1.50814 | 1.50818 | 1.50807 | 1.50811 |

The effect of finite proton size calculated here can be parametrized as $-0.0109 \langle r^2 \rangle$. However, higher iterations can change these results. For a very crude estimate, one can scale previous results for helium \cite{5} and assume that the ratio of nonperturbative to perturbative contributions was the same, giving a contribution of 0.175 meV. The contribution due to two and three iterations have been calculated by \cite{8} and \cite{23}, respectively, giving a total of 0.151 meV. An additional higher iteration including finite size and vacuum polarization is given in ref. \cite{8} (equations (66) and (67)) and ref. \cite{10} (equations (264) and (268)). These amount to $-0.0164 \langle r^2 \rangle$. The best way to calculate this would be an accurate numerical solution of the Dirac equation in the combined Coulomb-plus Uehling potential.

The mixed muon-electron vacuum polarization correction was recalculated and gave the same result as obtained previously, namely 0.00007 meV. \cite{21,10}.

The Wichmann-Kroll \cite{22} contribution was calculated using the parametrization for the potential given in \cite{11}. The result obtained (-0.00103 meV) is consistent with that given in \cite{10}, but not with that given in \cite{8}.

The equivalent potential for the virtual Delbrück effect was recomputed from the Fourier transform given in \cite{11} and \cite{1}. The resulting potential was checked by reproducing previously calculated results for the 2s-2p transition in muonic helium, and the 3d-2p transitions in muonic Mg and Si. The result for hydrogen is \((0.00135 \pm 0.00015)\) meV. As in the case of muonic helium, this contribution very nearly cancels the Wichmann-Kroll contribution. The contribution corresponding to three photons to the muon and one to the proton should be analogous to the light by light contribution to the muon anomalous moment; to my knowledge, the corresponding contribution to the muon form factor has never been calculated. It will be comparable to the other light by light contributions. For an estimate, the correction to the Lamb shift due to the contribution to the anomalous magnetic moment was calculated; it amounts to \((-)0.00002\) meV; the contribution to the muon form factor is one of the most significant unknown corrections.

The sixth order vacuum polarization corrections to the Lamb shift in muonic hydrogen have been calculated by Kinoshita and Nio \cite{23}. Their result for the 2p-2s transition is

$$
\Delta E^{(6)} = 0.120045 \cdot (\alpha Z)^2 \cdot m_e \left(\frac{\alpha}{\pi}\right)^3 \approx 0.00761 \text{meV}
$$

It is entirely possible that the as-yet uncalculated light by light contribution will give a comparable contribution.

The hadronic vacuum polarization contribution has been estimated by a number of authors \cite{24,25,10}. It amounts to about 0.012 meV. One point that should not be
forgotten about the hadronic VP correction is the fact that the sum rule or dispersion relation that everyone (including myself) used does not take into account the fact that the proton (nucleus) can in principle interact strongly with the hadrons in the virtual hadron loop. This is irrelevant for the anomalous magnetic moment but probably not for muonic atoms. An estimation of this effect appears to be extremely difficult, and could easily change the correction by up to 50%. Eides et al. [10] point out that the graph related to hadronic vacuum polarization can also contribute to the measured value of the nuclear charge distribution (and polarizability). It is not easy to determine where the contribution should be assigned.

Finite nuclear size and nuclear polarization

The main contribution due to finite nuclear size has been given analytically to order \((\alpha Z)^6\) by Friar [26]. The main result is

\[
\Delta E_{ns} = -\frac{2\alpha Z}{3} \left(\alpha Z m_r \right)^3 \left[ \langle r^2 \rangle - \frac{\alpha Z m_r}{2} \langle r^3 \rangle_2 \right] + (\alpha Z)^2 (F_{REL} + m_r^2 F_{NR}) \tag{5}
\]

where \(\langle r^2 \rangle\) is the mean square radius of the proton. For muonic hydrogen, the coefficient of \(\langle r^2 \rangle\) is \(5.1975\) (meV fm\(^{-2}\)), giving an energy shift (for the leading term) of \(3.862 \pm 0.108\) meV if the proton rms radius is \(0.862 \pm 0.012\) fm. The shift is \(4.163 \pm 0.188\) meV if the proton rms radius is \(0.895 \pm 0.018\) fm, and \(3.979 \pm 0.076\) meV if the proton rms radius of \(0.875 \pm 0.007\) fm. The second term in Eq. (5) contributes \(-0.0232\) meV for a dipole form factor and \(-0.0212\) meV for a Gaussian form factor. The parameters were fitted to the proton rms radius. This can be written as \(-0.0347\)\(\langle r^2 \rangle^{3/2}\) or \(0.0317\)\(\langle r^2 \rangle^{3/2}\), respectively. This differs slightly from the value given by Pachucki [9]. The model dependence introduces an uncertainty of about \(\pm 0.002\) meV. The remaining terms contribute \(0.00046\) meV. This estimate includes all of the terms given in [26], while other authors [9] give only some of them. Clearly the neglected terms are not negligible. There is also a contribution of \(-3 \cdot 10^{-6}\) meV to the binding energy of the 2\(p_{1/2}\)-level, and a recoil correction of \(0.012\) meV to the binding energy of the 2s-level.

As mentioned previously, the finite-size contributions to vacuum polarization can be parametrized as \(-0.0109\langle r^2 \rangle - 0.0164\langle r^2 \rangle\), giving a total of \(-0.0273\langle r^2 \rangle\) or \(-0.0209(6)\) meV if the proton radius is \(0.875\) fm.

The contribution due to nuclear polarization has been calculated by Rosenfelder [27] to be \(0.017 \pm 0.004\) meV, and by Pachuki [9] to be \(0.012 \pm 0.002\) meV. Other calculations [28, 29] give intermediate values (0.013 meV and 0.016 meV, respectively). The value appearing in table 1 is an average of the three most recent values, with the largest quoted uncertainty, which is probably underestimated.

Relativistic Recoil

As is well-known, the center-of-mass motion can be separated exactly from the relative motion only in the nonrelativistic limit. Relativistic corrections have been studied by many authors, and will not be reviewed here. The relativistic recoil corrections summarized in [11] include the effect of finite nuclear size to leading order in \(m_\mu/m_N\) properly.
Up to now this method has been used to treat recoil corrections to vacuum polarization only in the context of extensive numerical calculations that include the Uehling potential in the complete potential, as described in [1]. They can be included explicitly, as a perturbation correction to point-Coulomb values. Recall that (to leading order in $1/m_N$), the energy levels are given by

$$E = E_r - \frac{B_0^2}{2m_N} + \frac{1}{2m_N} \langle h(r) + 2B_0P_1(r) \rangle$$

(6)

where $E_r$ is the energy level calculated using the reduced mass and $B_0$ is the unperturbed binding energy. Also

$$h(r) = -P_1(r)[P_1(r) + \frac{1}{r}Q_2(r)] - \frac{1}{3r}Q_2(r)[P_1(r) + Q_4(r)/r^3]$$

(7)

Here

$$P_1(r) = 4\pi\alpha Z \int_r^{\infty} r'\rho(r')dr' = -V(r) - rV'(r)$$

(8)

$$Q_2(r) = 4\pi\alpha Z \int_0^r r^2\rho(r')dr' = r^2V'(r)$$

$$Q_4(r) = 4\pi\alpha Z \int_0^r r^4\rho(r')dr'$$

An effective charge density $\rho_{VP}$ for vacuum polarization can be derived from the Fourier transform of the Uehling potential. Recall that (for a point nucleus)

$$V_{Uehl}(r) = -\frac{\alpha Z}{r} \frac{2\alpha}{3\pi} \cdot \chi_1(2m_e r)$$

$$= -(\alpha Z) \frac{2\alpha}{3\pi} \cdot \int_1^{\infty} dz \frac{(z^2 - 1)^{1/2}}{z^2} \cdot \left(1 + \frac{1}{2z^2}\right) \left( \frac{2}{\pi} \int_0^{\infty} \frac{q^2 \cdot j_0(qr)}{q^2 + 4m_e^2 z^2} dq \right)$$

(9)

where $\chi_n(x)$ is defined in [1]. In momentum space, the Fourier transform of $\nabla^2 V$ is obtained by multiplying the Fourier transform of $V$ by $-q^2$. Note that using the normalizations of [1, 6], one has $\nabla^2 V = -4\pi\alpha Z \rho$ where $\rho$ is the charge density. One then obtains

$$4\pi \rho_{VP}(r) = \frac{2\alpha}{3\pi} \cdot \int_1^{\infty} dz \frac{(z^2 - 1)^{1/2}}{z^2} \cdot \left(1 + \frac{1}{2z^2}\right) \left( \frac{2}{\pi} \int_0^{\infty} \frac{q^4 \cdot j_0(qr)}{q^2 + 4m_e^2 z^2} dq \right) = \frac{2}{\pi} \cdot \int_0^{\infty} q^2 U_2(q) j_0(qr) dq$$

(9)

$U_2(q)$ is defined in [1]. It is also easy to show that

$$\frac{dV_{Uehl}}{dr} = +\frac{\alpha Z}{r} \frac{2\alpha}{3\pi} \cdot \left[ \frac{1}{r} \chi_1(2m_e r) + 2m_e \chi_0(2m_e r) \right]$$

$$= -\frac{1}{r} V_{Uehl}(r) + (\alpha Z) \frac{2\alpha}{3\pi} \cdot \frac{2m_e}{r} \chi_0(2m_e r)$$
Keeping only the Coulomb and Uehling potentials, one finds

$$P_1(r) = -\alpha Z \frac{2\alpha}{3\pi} (2m_e) \chi_0 (2m_e r)$$

$$Q_2(r) = \alpha Z \left( 1 + \frac{2\alpha}{3\pi} [\chi_1 (2m_e r) + (2m_e r) \chi_0 (2m_e r)] \right)$$

$$Q_4(r) = \alpha Z \frac{2\alpha}{3\pi} \int_1^\infty dz \frac{(z^2 - 1)^{1/2}}{z^2} \left( 1 + \frac{1}{2z^2} \right) \cdot \left( \frac{2}{\pi} \right) \int_0^\infty \frac{1}{q^2 + 4m_e^2 z^2} \frac{[6qr - (qr)^3] \cos(qr) + [3(qr)^2 - 6] \sin(qr)}{q} dq$$

where $\chi_n(x)$ is defined in [1]. Corrections due to finite nuclear size can be included when a model for the charge distribution is given. This done by Friar [26] (and confirmed independently for two different model charge distributions); the contribution due to finite nuclear size to the recoil correction for the binding energy of the 2s-level is -0.013 meV. The factor $1/m_n$ is replaced by $1/(m_\mu + m_N)$, also consistent with the calculations presented in [26].

Since vacuum polarization is assumed to be a relatively small correction to the Coulomb potential, it will be sufficient to approximate $Q_2(r)$ by $\alpha Z/r$. After some algebra, one can reduce the expectation values to single integrals:

$$\langle P_1(r) \rangle = 2m_e \alpha Z \frac{2\alpha}{3\pi} \int_1^\infty \frac{(z^2 - 1)^{1/2}}{z} \cdot \left( 1 + \frac{1}{2z^2} \right) \cdot \left( \frac{(az)^2 - az + 1}{(1 + az)^5} \delta_{\ell_0} + \frac{1}{(1 + az)^5} \delta_{\ell_1} \right) dz \quad (10)$$

$$\left\langle \frac{\alpha Z}{r} P_1(r) \right\rangle = - (\alpha Z)^3 m_r m_e \frac{2\alpha}{3\pi} \int_1^\infty \frac{(z^2 - 1)^{1/2}}{z} \cdot \left( 1 + \frac{1}{2z^2} \right) \cdot \left( \frac{2(az)^2 + 1}{2(1 + az)^4} \delta_{\ell_0} + \frac{1}{2(1 + az)^4} \delta_{\ell_1} \right) dz \quad (11)$$

with $a = 2m_e/(\alpha Z m_r)$. When Eq. (10) is multiplied by $-2B_0/(m_\mu + m_N)$ this results in a shift of $-0.00015$ meV for the 2s-state and of $-0.0001$ meV for the 2p-state, and when Eq. (11) is multiplied by $1/(m_\mu + m_N)$ this results in a shift of 0.00489 meV for the 2s-state and of 0.0017 meV for the 2p-state. These expectation values also appear when vacuum polarization is included in the Breit equation [31].

Finally,

$$\langle \frac{\alpha Z}{3r^4} Q_4(r) \rangle = - (\alpha Z)^4 \frac{m_r^2}{6} \frac{2\alpha}{3\pi} \int_1^\infty \frac{(z^2 - 1)^{1/2}}{z^2} \cdot \left( 1 + \frac{1}{2z^2} \right) \cdot \left[ - \frac{6}{az} \left( \frac{2 + az}{1 + az} - \frac{2}{az} \ln(1 + az) \right) + \frac{3(az)^2 + 2az - 1}{(1 + az)^3} + \frac{3 + az}{4(1 + az)^4} \delta_{\ell_0} + \frac{1 - 3az - 2(az)^2}{4(1 + az)^4} \delta_{\ell_1} \right] dz \quad (12)$$
When multiplied by $1/(m_\mu + m_N)$ this results in a shift of 0.002475 meV for the 2s-state and of 0.000238 meV for the 2p-state.

Combining these expectation values according to equations \[3\] and \[4\] one finds a contribution to the 2p-2s transition of -0.00419 meV. To obtain the full relativistic and recoil corrections, one must add the difference between the expectation values of the Uehling potential calculated with relativistic and nonrelativistic wave functions, giving a total correction of 0.0166 meV. This is in fairly good agreement with the correction of 0.0169 meV calculated by Veitia and Pachucki \[31\], using a generalization of the Breit equation \[32\] which is similar to that given in \[6\]. The treatment presented here has the advantage of avoiding second order perturbation theory.

The review by Eides et.al \[10\] gives a better version of the two photon recoil (Eq. 136) than was available for the review by Borie and G. Rinker \[1\]. Evaluating this expression for muonic hydrogen gives a contribution of -0.04497 meV to the 2p-2s transition. Higher order radiative recoil corrections give an additional contribution of -0.0096 meV \[10\]. However, some of the contributions to the expressions given in \[10\] involve logarithms of the mass ratio $m_\mu/m_N$. Logarithms can only arise in integrations in the region from $m_\mu$ to $m_N$; in this region the effect of the nuclear form factor should not be neglected. Pachucki \[8\] has estimated a finite size correction to this of about 0.02 meV, which seems to be similar to the term proportional to $\langle r^3 \rangle_{(2)}$ given in Eq.\[5\] as calculated in the external field approximation by Friar \[26\]. This two-photon correction requires further investigation. In particular, the parametrization of the form factors used in any calculation should reproduce the correct proton radius.

An additional recoil correction for states with $\ell \neq 0$ has been given by \[32\] (see also \[10\]). It is

$$\Delta E_{n,\ell,j} = \left(\frac{\alpha Z}{2n^3 m_N^2}\right)^4 \cdot \frac{m_\mu^3}{2n^3 m_N^2} \cdot \left(1 - \delta_{\ell 0}\right) \left(1 - \frac{1}{\kappa(2\ell + 1)}\right)$$

(13)

When evaluated for the 2p-states of muonic hydrogen, one finds a contribution to the 2p-2s transition energy of 0.0575 meV for the 2p\(_{1/2}\) state and -0.0287 meV for the 2p\(_{3/2}\) state.

**Muon Lamb Shift**

For the calculation of muon self-energy and vacuum polarization, the lowest order (one-loop approximation) contribution is well-known, at least in perturbation theory. Including also muon vacuum polarization (0.0168 meV) and an extra term of order $(Z\alpha)^5$ as given in \[10\]:

$$\Delta E_{2s} = \frac{\alpha (\alpha Z)^5 m_\mu}{4} \cdot \left(\frac{m_r}{m_\mu}\right)^3 \cdot \left(\frac{139}{64} + \frac{5}{96} - \ln(2)\right)$$

which contributes -0.00443 meV, one finds a contribution of -0.66788 meV for the 2s\(_{1/2} - 2p_{1/2}\) transition and -0.65031 meV for the 2s\(_{1/2} - 2p_{3/2}\) transition.

A misprint in the evaluation of the contribution of the higher order muon form factors (contributing to the fourth order terms) has been corrected. The extra electron loop contribution to $F_2(0)$ is should be $1.09426 (\alpha/\pi)^2$. This reproduces the correct coefficient of $(\alpha/\pi)^2$ from the muon $(g-2)$ analyses. This is .7658, which is equal to 1.09426-0.32848.
The fourth order electron loops \[30\] dominate the fourth order contribution (-0.00169 meV and -0.00164 meV, respectively). The rest is the same as for the electron \[1\]. The contribution of the electron loops alone is -0.00168 meV for the \(2s_{1/2} - 2p_{1/2}\) transition and -0.00159 meV for the \(2s_{1/2} - 2p_{3/2}\) transition.

Pachuki \[8\] has estimated an additional contribution of -0.005 meV for a contribution corresponding to a vacuum polarization insert in the external photon.

**Summary of contributions**

Using the fundamental constants from the CODATA 1998 \([12]\) one finds the transition energies in meV in table 1. Here the main vacuum polarization contributions are given for a point nucleus, using the Dirac equation with reduced mass. Some uncertainties have been increased from the values given by the authors, as discussed in the text.

The finite size corrections up to order \((\alpha Z)^5\) can be parametrized as

\[
5.1975\langle r^2 \rangle - 0.0109\langle r^2 \rangle - 0.0164\langle r^2 \rangle + 0.0347\langle r^3 \rangle.
\]

In the case of the muon Lamb shift, the numbers in table 1 are for the \(2s_{1/2} - 2p_{1/2}\) transition. The corresponding numbers for the \(2s_{1/2} - 2p_{3/2}\) transition are -0.65031 meV and -0.00164 meV, respectively.

**Fine structure of the 2p state**

There are two possible ways to calculate the fine structure. One is to start with the point Dirac value, include the contribution due to vacuum polarization, as calculated above, as well as the spin-orbit splitting (computed perturbatively) due to the muon’s anomalous magnetic moment, and recoil as given by Eq. \([13]\). The results are summarized in table 2.

An alternative method is to use the formalism given in \([6]\) (and elsewhere, see, eg. \([32, 10]\)) which gives the energy shift as the expectation value of

\[
-\frac{1}{r} \frac{dV}{dr} \cdot \frac{1 + a_\mu + (a_\mu + 1/2)m_N/m_\mu \vec{L} \cdot \vec{\sigma}_\mu}{m_N m_\mu}.
\]

(14)

Note that

\[
\frac{1}{m_N m_\mu} + \frac{1}{2m_\mu^2} = \frac{1}{2m_r^2} - \frac{1}{2m_N^2}
\]

so that the terms not involving \(a_\mu\) in the spin-orbit contribution are really the Dirac fine structure plus the Barker-Glover correction (Eq. \[13\]).

The Uehling potential has to be included in the potential \(V(r)\). For states with \(\ell > 0\) in light atoms, and neglecting the effect of finite nuclear size, we may take

\[
\frac{1}{r} \frac{dV}{dr} = \frac{\alpha Z}{r^3} \cdot \left[ 1 + \frac{2\alpha}{3\pi} \int_1^\infty \frac{(z^2 - 1)^{1/2}}{z^2} \cdot \left( 1 + \frac{1}{2z^2} \right) \cdot (1 + 2m_\mu rz) \cdot e^{-2m_\mu rz} dz \right]
\]

(15)

which is obtained from the Uehling potential \([13, 14]\) by differentiation. Then, assuming that it is sufficient to use nonrelativistic point Coulomb wave functions for the 2p state, one finds

\[
\langle \frac{1}{r^3} \rangle_{2p} \rightarrow \langle \frac{1}{r^3} \rangle_{2p} \cdot (1 + \varepsilon_{2p})
\]
| Contribution                          | Value (meV) | Uncertainty (meV) |
|--------------------------------------|-------------|-------------------|
| Uehling                              | 205.0282    |                   |
| Källen-Sabry                         | 1.5081      |                   |
| Wichmann-Kroll                       | -0.00103    |                   |
| virt. Delbrueck                      | 0.00135     | 0.00015           |
| mixed mu-e VP                        | 0.00007     |                   |
| hadronic VP                          | 0.011       | 0.002             |
| sixth order [23]                     | 0.00761     |                   |
| recoil [10] (eq136)                  | -0.04497    |                   |
| recoil, higher order [10]            | -0.0096     |                   |
| recoil, finite size [26]             | 0.013       | 0.001             |
| recoil correction to VP [1]          | -0.0041     |                   |
| additional recoil [32]               | 0.0575      |                   |

| muon Lamb shift                      |             |                   |
| second order                         | -0.66788    |                   |
| fourth order                         | -0.00169    |                   |
| nuclear size ($R_p=0.875$ fm)        | 0.007 fm    |                   |
| main correction [26]                 | -3.979      | 0.076             |
| order ($\alpha Z$)$^5$ [26]          | 0.0232      | 0.002             |
| order ($\alpha Z$)$^6$ [26]          | -0.0005     |                   |
| correction to VP                     | -0.0083     |                   |
| polarization                         | 0.015       | 0.004             |
| Other (not checked)                  |             |                   |
| VP iterations [8]                    | 0.151       |                   |
| VP insertion in self energy [8]      | -0.005      |                   |
| additional size for VP [10]          | -0.0128     |                   |

Table 1: Contributions to the muonic hydrogen Lamb shift. The proton radius is taken from [20].

| $E(2p_{3/2}) - E(2p_{1/2})$ (meV)          |
|---------------------------------------------|
| Dirac                                       | 8.41564   |
| Uehling(VP)                                 | 0.0050    |
| Källen-Sabry                                | 0.00004   |
| anomalous moment $a_\mu$                    |           |
| second order                                | 0.01757   |
| higher orders                               | 0.00007   |
| Recoil (Eq.(13))                            | -0.0862   |
| Total Fine Structure                         | 8.352     |

Table 2: Contributions to the fine structure of the 2p-state in muonic hydrogen.
where
\[ \varepsilon_{2p} = \frac{2\alpha}{3\pi} \int_{1}^{\infty} \frac{(z^2 - 1)^{1/2}}{z^2} \cdot \left( 1 + \frac{1}{2z^2} \right) \cdot \left( 1 + \frac{2az}{(1 + az)^3} \right) \, dz \]  
(16)
with \( a = 2m_e/(\alpha Zm_r) \). The result for the fine structure is
\[ \frac{-(\alpha Z)^4m_r^3}{n^3(2\ell + 1)\kappa} \cdot \left( \frac{1}{m_Nm_\mu} + \frac{1}{2m_\mu^2} + \frac{a_\mu}{m_\mu m_r} \right) \cdot (1 + \varepsilon_{2p}) \]  
(17)
where \( \varepsilon_{2p} \) is given by Eq. (16). In this case, the terms involving \( a_\mu \) in the expression for the muon Lamb shift are included, and should not be double counted. With a numerical value of \( \varepsilon_{2p} = 0.000365 \), one finds a contribution of 0.00305 meV (compared with 0.005 meV using Dirac wave functions).

Numerically, the terms not involving \( a_\mu \) give a contribution of 8.3291 meV and the contribution from \( a_\mu \) gives a contribution of 0.0176 meV, for a total of 8.3467 meV, in good agreement with Eq. 80 of [8]. When the vacuum polarization correction is added, the result is only very slightly different from the Dirac value of 8.352 meV. The contribution due to the anomalous magnetic moment of the muon is the same in both cases.

In both cases one should include the \( B^2/2M_N \)-type correction to the fine structure. (see [10], Eq(38)). This is tiny \((5.7 \cdot 10^{-6} \text{ meV})\) and is not included in the table. Friar [26] has given expressions for the energy shifts of the 2p-states due to finite nuclear size. These were calculated and found to give a negligible contribution \((3.1 \cdot 10^{-6} \text{ meV})\) to the fine structure of the 2p-state.

**Hyperfine structure**

The hyperfine structure is calculated in the same way as was done in earlier work [6, 7], but with improved accuracy. Most of the formalism and results are similar to those given by [8].

**The 2p state:**
The hyperfine structure of the 2p-state is given by \[ \hat{\mathcal{H}}_{\text{HFS}} = \hat{\mathcal{H}}_{\text{spin}} + \hat{\mathcal{H}}_{\text{iso}} \]  
(18)
where \( \hat{j} = \sqrt{2j + 1} \), the 6-j symbols are defined in [33], and \( C_{F1} = \delta_{F1} - 2\delta_{F0} - (1/5)\delta_{F2} \).

\[ x = \frac{m_\mu(1 + 2\kappa_p)}{2m_N(1 + \kappa_p)} \]
represents a recoil correction due to Thomas precession [6, 32]. The same correction due to vacuum polarization (Eq. (10)) should be applied to the HFS shifts of the 2p-states, as well as to the spin-orbit term.
As has been known for a long time [6, 7, 8], the states with total angular momentum \( F = 1 \) are a superposition of the states with \( j = 1/2 \) and \( j = 3/2 \). Let the fine structure splitting be denoted by \( \delta = E_{2p3/2} - E_{2p1/2} \), and let

\[
\beta = \frac{(\alpha Z)^4 m_\mu^3}{3m_\mu m_N} \cdot (1 + \kappa_p)
\]

and \( \beta' = \beta \cdot (1 + \varepsilon_{2p}) \).

The energy shifts of the 2p-states with total angular momentum \( F \) (notation \( ^{2F+1}L_j \)) are then given in Table 3:

| State   | Energy | Energy in meV |
|---------|--------|---------------|
| \(^1p_{1/2}\) | \(-\beta'(2 + x + a_\mu)/8\) | -5.971 |
| \(^3p_{1/2}\) | \((\Delta - R)/2\) | 1.846 |
| \(^3p_{3/2}\) | \((\Delta + R)/2\) | 6.376 |
| \(^5p_{3/2}\) | \(\delta + \beta'(1 + 5x/4 - a_\mu/4)/20\) | 9.624 |

Table 3: Hyperfine structure of the 2p-state in muonic hydrogen.

where

\[
\Delta = \delta - \beta'(x - a_\mu)/16
\]

\[
R^2 = [\delta - \beta'(1 + 7x/8 + a_\mu/8)/6]^2 + (\beta')^2(1 + 2x - a_\mu)^2/288
\]

(Here \( \delta = 8.352 \text{ meV} \) Some minor errors in [6] have been corrected. These numbers differ slightly from those given in ref. [10].

The 2s-state:
The basic hyperfine splitting of the 2s-state is given by

\[
\Delta \nu_F = \frac{(\alpha Z)^4 m_\mu^3}{3m_\mu m_N} \cdot (1 + \kappa_p) \cdot (1 + a_\mu) = \beta \cdot (1 + a_\mu) = 22.8332 \text{ meV}
\]

(see, for example [10], Eq. (271,277)) As was shown in [6, 10], the energy shift of the 2s-state is given by:

\[
\Delta E_{2s} = \beta \cdot (1 + a_\mu) \cdot (1 + \varepsilon_{VP} + \varepsilon_{\text{vertex}} + \varepsilon_{\text{Breit}} + \varepsilon_{FS,rec}) \cdot [\delta_{F1} - 3\delta_{F0}] / 4
\]

Here (34)

\[
\varepsilon_{\text{vertex}} = \frac{2\alpha (\alpha Z)}{3} \left( \ln(2) - \frac{13}{4} \right) = -1.36 \cdot 10^{-4}
\]

and (10, Eq. (277))

\[
\varepsilon_{\text{Breit}} = \frac{17(\alpha Z)^2}{8} = 1.13 \cdot 10^{-4}
\]
The vacuum polarization correction has two contributions. One of these is a result of a modification of the magnetic interaction between the muon and the nucleus and is given by (see [7])

\[
\varepsilon_{VP1} = \frac{4\alpha}{3\pi^2} \int_0^\infty r^2 \, dr \left( \frac{R_{ns}(r)}{R_{ns}(0)} \right)^2 \int_0^\infty q^4 j_0(qr) G_M(q) \, dq \\
\int_1^\infty \frac{(z^2 - 1)^{1/2}}{z^2} \left( 1 + \frac{1}{2z^2} \right) \frac{dz}{4m_e^2 z^2 + (q/2m_e)^2}
\]

(20)

One can do two of the integrals analytically and obtains for the 2s-state (with \( a = 2m_e/(\alpha Z m_r) \) and \( \sinh(\phi) = q/(2m_e) = K/a \))

\[
\varepsilon_{VP1} = \frac{4\alpha}{3\pi^2} \int_0^\infty \frac{K^2}{(1 + K^2)^2} F(\phi) G_M(\alpha Z m_r K) \, dK \left[ 2 - \frac{7}{(1 + K^2)^2} + \frac{6}{(1 + K^2)^2} \right]
\]

(21)

where \( F(\phi) \) is known from the Fourier transform of the Uehling potential and is given by Eq.(4).

The other contribution, as discussed by [34, 35] arises from the fact that the lower energy hyperfine state, being more tightly bound, has a higher probability of being in a region where vacuum polarization is large. This results in an additional energy shift of

\[
2 \int V_{Uehl}(r) \psi_{2s}(r) \delta M \psi_{2s}(r) d^3r
\]

Following Ref. [34] with \( y = (\alpha Z m_r/2) \cdot r \), one has

\[
\delta M \psi_{2s}(r) = 2m_\mu \Delta \nu_F \psi_{2s}(0) \left( \frac{2}{\alpha Z m_r} \right)^2 \exp(-y) \left[ (1 - y)(\ln(2y) + \gamma) + \frac{13y - 3 - 2y^2}{4} - \frac{1}{4y} \right]
\]

(\( \gamma \) is Euler’s constant), and

\[
\psi_{2s}(r) = \psi_{2s}(0)(1 - y) \exp(-y)
\]

One finds after a lengthy integration

\[
\varepsilon_{VP2} = \frac{16\alpha}{3\pi^2} \int_0^\infty \frac{dK \, G_E(\alpha Z m_r K) F(\phi)}{1 + K^2} \left\{ \frac{1}{2} - \frac{17}{(1 + K^2)^2} + \frac{41}{(1 + K^2)^3} - \frac{24}{(1 + K^2)^4} + \frac{\ln(1 + K^2)}{1 + K^2} \left[ 2 - \frac{7}{(1 + K^2)^2} + \frac{6}{(1 + K^2)^2} \right] + \frac{\tan^{-1}(K)}{K} \left[ 1 - \frac{19}{2(1 + K^2)^2} + \frac{20}{(1 + K^2)^2} - \frac{12}{(1 + K^2)^3} \right] \right\}
\]

(22)

Sternheim [35] denotes the two contributions by \( \delta_M \) and \( \delta_E \), respectively. An alternative expression, obtained by assuming a point nucleus, using Eq.(131) from [1] for the Uehling
potential, and doing the integrations in a different order, is
\[
\varepsilon_{VP2} = \frac{16\alpha}{3\pi} \int_1^{\infty} \frac{\left(z^2 - 1\right)^{1/2}}{z^2} \cdot \left(1 + \frac{1}{2z^2}\right) \cdot \frac{1}{(1 + az)^2} \\
\times \left[\frac{az}{2} - \frac{1}{1+az} + \frac{23}{8(1 + az)^2} - \frac{3}{2(1 + az)^3}\right] dz
\]
with \(a = 2m_e/(\alpha Z m_{\text{red}})\). Both methods give the same result.

In the case of ordinary hydrogen, each of these contributes \(3\alpha^2/8 = 1.997 \cdot 10^{-5}\). The accuracy of the numerical integration was checked by reproducing these results. One can thus expect that muonic vacuum polarization will contribute \(3\alpha^2/4 \simeq 4 \cdot 10^{-5}\), as in the case of normal hydrogen. This amounts to an energy shift of 0.0009 meV. Contributions due to the weak interaction or hadronic vacuum polarization should be even smaller.

For muonic hydrogen, one obtains \(\varepsilon_{VP1} = 0.00211\) and \(\varepsilon_{VP2} = 0.00325\) for a point nucleus. Including the effect of the proton size (with \(G_E(q) = G_M(q)\) as a dipole form factor) reduces these numbers to 0.00206 and 0.00321, respectively. For the case of muonic \(^3\)He, the corresponding numbers are \(\varepsilon_{VP1} = 0.00286\) and \(\varepsilon_{VP2} = 0.00476\). The contribution to the hyperfine splitting of the 2s-state is then 0.0470 meV+0.0733 meV=0.1203 meV (0.1212 meV if muonic vacuum polarization is included). The combined Breit and vertex corrections reduce this value to 0.1207 meV. (0.1226 meV if the proton form factors are not taken into account).

The correction due to finite size and recoil have been given in \[8\] as -0.145 meV, while a value of -0.152 meV is given in \[38\]. Ref. \[8\] also gives a correction as calculated by Zemach (\[36\]) equal to -0.183 meV, but claims that this correction does not treat recoil properly. The Zemach correction is equal to
\[
\varepsilon_{Zem} = -2\alpha Z m_r \langle r \rangle_{(2)}
\]
where \(\langle r \rangle_{(2)}\) is given in \[31\] \[29\] \[37\]. Using the value \(\langle r \rangle_{(2)} = 1.086 \pm 0.012 \text{ fm}\) from \[37\], gives \(\varepsilon_{Zem} = -0.00702\), and a contribution of -0.1742 meV to the hyperfine splitting of the 2s state. Including this, but not other recoil corrections to the hyperfine structure of the 2s-state gives a total splitting of 22.7806 meV. Additional higher order corrections calculated in Ref. \[38\] amount to a total of -0.0003 meV and are not included here.

**Summary of contributions and Conclusions**

The most important contributions to the Lamb shift in muonic hydrogen, including hyperfine structure, have been independently recalculated. A new calculation of some terms that were omitted in the most recent literature, such as the virtual Delbrück effect \[11\] and an alternative calculation of the relativistic recoil correction have been presented.
Numerically the results given in table 1 add up to a total correction of
(206.032(6) − 5.225 ⟨r^2⟩ + 0.0347 ⟨r^2⟩^{3/2}) \text{meV} = 202.055\pm0.12\text{meV}. (for the value of the
proton radius from [20]). As is well known, most of the uncertainty arises from the
uncertainty in the proton radius.

However, the contribution of the light-by-light graph to the muon form factor has not
yet been calculated. Also, since \(m_\mu/m_p = 0.1126\) is much larger than \(\alpha Z\), it is possible
that recoil corrections of higher order in the mass ratio, that have never been calculated,
could be significant at the level of the expected experimental accuracy of about 0.01 meV.
In particular, the two-photon recoil corrections, including finite nuclear size, should be
recalculated to resolve (small) inconsistencies among various theoretical results.

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References
[1] E. Borie, G.A. Rinker, Rev. Mod. Phys. 54, 67 (1982)
[2] F. Kottmann et al., Hyperf. Int. 138, 55 (2001)
[3] A. di Giacomo, Nucl. Phys. B11, 411 (1969)
[4] E. Borie, Z. Phys. A275, 347 (1975)
[5] E. Borie, G.A. Rinker, Phys. Rev. A18, 324 (1978)
[6] E. Borie, Z. Phys. A278, 127 (1976)
[7] E. Borie, Z. Phys. A297, 17 (1980)
[8] K. Pachucki, Phys. Rev. A53, 2092 (1996)
[9] K. Pachucki, Phys. Rev. A60, 3593 (1999)
[10] M.I. Eides, H. Grotch, V.A. Shelyuto, Physic Reports, 342, 63-261 (2001)
[11] E. Borie, Nucl. Phys. A267, 485 (1976)
[12] P.J. Mohr, B.N.Taylor, Rev. Mod. Phys. 72, 351 (2000)
[13] E.A. Uehling, Phys. Rev. 48, 55 (1935)
[14] R. Serber, Phys. Rev. 48, 49 (1935)
[15] A.I. Akhiezer, V.B. Berestetskii, Quantum electrodynamics, Wiley Interscience, New York,
1965.
[16] G. Källen, A. Sabry, K. Dan. Vidensk. Selsk. Mat.-Fys. Medd. 29, #17 (1955)
[17] G.G. Simon et al., Nucl. Phys. A333, 38 (1980)
[18] R. Rosenfelder, Phys. Lett. B479, 381 (2000)
[19] I. Sick, Phys. Lett. B576, 62 (2003)
[20] P.J. Mohr, B.N. Taylor, Rev. Mod. Phys. 76, (2004), in press
[21] E. Borie, Helv. Phys. Acta 48, 671 (1975)
[22] E.H. Wichmann, N.M. Kroll, Phys. Rev. 101, 843 (1956)
[23] T. Kinoshita, M. Nio, Phys. Rev. Lett. 82, 3240 (1999)
[24] E. Borie, Z. Phys. A302, 187 (1981)
[25] J.L. Friar, J. Martorell, D.W.L. Sprung, Phys. Rev. A59, 4061 (1999)
[26] J.L. Friar, Annals of Physics, 122, 151 (1979)
[27] R. Rosenfelder, Phys. Lett. B463, 317 (1999)
[28] S.A. Srartsev, et al., Phys. Atom Nucl. 1233 (1976)
[29] R.N. Faustov, A.P. Martynenko, AIP Conf. Proc. 564, 277 (2001)
[30] R. Barbieri, M. Caffo, E. Remiddi, Lett. Nuov. Cim. 7, 60 (1973)
[31] A. Veitia, K. Pachucki, Phys. Rev. A69, 042501 (2004)
[32] E.H. Barker, N.M. Glover, Phys. Rev. 99, 317 (1955)
[33] A.R. Edmonds, *Angular Momentum in Quantum Mechanics*, Princeton University Press, 1960
[34] S.J. Brodsky, G.W. Erickson, Phys. Rev. 148, 26 (1966)
[35] M.M. Sternheim, Phys. Rev. 138, B430 (1965)
[36] A.C. Zemach, Phys. Rev. 104, 1771 (1956)
[37] J.L. Friar, I. Sick, Phys. Lett. B579, 285 (2004)
[38] A.P. Martynenko, Preprint SSU-HEP-04/09