Analysis and Performance of Lumped-Element Kinetic Inductance Detectors for W-band

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Abstract—Lumped-Element Superconducting resonators are a promising technology for its use in millimeter-wave observations and quantum computing applications that require large arrays of extremely sensitive detectors. Among them, Lumped-Element Kinetic Inductor Detectors (LEKIDs) have shown good performance in the submillimeter band in several Earth-based telescopes. In this work, LEKIDs for their use as millimeter-wave receivers of astronomical applications are presented. LEKIDs arrays using a thin bilayer of superconducting titanium/aluminum (Ti/Al), deposited on silicon substrate, have been designed and fabricated. The design of a dual-polarization LEKID with the goal of detection at W-Band for two orthogonal polarizations is described and a fabricated array has demonstrated absorption at ambient temperature. Also, an approximate design methodology of the coupling parameter for LEKIDs readout, essential for dynamic range optimization of the detector under millimeter wave radiation, is proposed. In addition, resonance characteristics and coupling factor of the fabricated superconducting resonators using high-quality internal factor Q under cryogenic temperatures have been analyzed. The design guidelines in this work are applicable to other LEKIDs arrays, and the presented superconducting Ti/Al thin film LEKIDs can be used in future receiver arrays in millimeter bands.

Index Terms— Kinetic Inductance Detector, superconducting microwave devices, lumped-element resonator, cryogenics, millimeter wave astronomy, polarimetry.

I. INTRODUCTION

KINETIC Inductance Detectors (KIDs) are used in high-sensitivity radio astronomy receivers. These detectors operate under the property of inductance change of a superconducting strip, when it is cooled down well below its superconducting critical temperature, usually near absolute zero, when they absorb a microwave signal taken by a telescope receiver system. KIDs based receivers are built containing thousands of individual detectors, or instrument pixels, as their intrinsic multiplexing is a relatively simple way of receiver integration. Frequency domain multiplex operation consists of a high amount of resonant circuits coupled to a single readout transmission line, making feasible such a high number of instrument pixels relatively easy [1]. Each individual KID is embedded in a resonant circuit having a resonant frequency different from the remainder KIDs. In this way, the readout system is in fact a frequency division multiplex communication system [2]. Former KID receivers operated in submillimeter bands, but, nowadays, the use of new superconducting materials, allows detection in the optical-near IR band [3] or even in the millimeter band [4],[5] showing good performance in all cases.

KIDs are based on superconducting microresonators with a characteristic resonant frequency and quality factor. The absorption of an incoming radiation, with energy higher than the superconducting gap, changes the kinetic inductance (L_k), leading to a change of the resonant frequency. Moreover, this absorption also produces dissipative losses in the resonator, decreasing the quality factor in comparison with the darkness condition. The reduction of resonant frequency and quality factor, given proper calibration, enables the detection of radio signals from astronomical sources, employing complex processing algorithms.

Several KID radio astronomy instruments for millimeter and submillimeter wave observations are in operation, achieving even thousands of multiplexed detectors [6]-[10]. There are two main types of detectors: Antenna coupled- Microwave Kinetic Inductance Detectors (MKIDs) [2] and Lumped Element Kinetic Inductance Detectors (LEKIDs) [11]. The former are based on distributed resonators with coplanar-waveguide (CPW) geometries coupled to an antenna that acts as an active receiver. The later, LEKIDs, which are the focus of our study, are based on lumped inductor-capacitor resonators, where the inductor acts directly as the effective optical absorber, and its design is shown in Fig. 1.
Fig. 1. Structure of a LEKID (Lumped Kinetic Inductance Detector): Meander Inductor with strips of width $2a$ and spacing $g$. Coupling separation $s$ between the resonator and the readout 50 Ohm microstrip feedline of width $W_1$.

Great efforts have been made to develop dual-polarization LEKIDs and to improve the sensitivity at W-band for future Cosmic Microwave Background observations. For instance, A. Catalano et al. [5], [12], [13], showed that the superconducting proximity effect in a titanium/aluminum (Ti/Al) bilayer causes a decrease in the superconducting critical temperature that pushes the LEKID detection to the W-band. Regarding polarization sensitivity, earth-based instruments such as NIKA2 [9] use an external polarizer that separates the two linear polarizations into two independent arrays based on a Hilbert fractal structure. On the other hand, some additional works have focused on the development of dual-polarization LEKIDs demonstrating simultaneous orthogonal polarization sensitivity in the millimeter and submillimeter wave ranges [8], [14]. However, very little has been done in developing dual-polarization LEKIDs in the W-band. Therefore, in this work, we focus on both requirements and present an optimized version of the BiKID approach [15] for the W-band, based on an optimized optical coupling design.

The response of LEKIDs is maximized when critical coupling is achieved, i.e. the external quality factor ($Q_c$) equals the internal quality factor ($Q_i$) [16]. Whereas $Q_i$ is set by fixed parameters such as the optical background or operating base temperature, $Q_c$ can be modified by tuning geometrical parameters as separation $s$. Therefore, a design methodology for the microwave coupling to the LEKID has been developed and applied to a new prototype, with several external factors $Q_c$.

This paper presents the analysis, design, fabrication and experimental tests of bilayer LEKID prototypes working in W-band (75 to 110 GHz). The remainder of this article is organized as follows. Section II deals with the KID design, concerning millimeter wave signals to be detected, and low frequency resonators design, with special focus on external coupling. Section III describes fabrication and prototypes assembly. Section IV presents experimental systems to perform ambient temperature and cryogenic tests. Tests results are detailed in Section V. Finally, Section VI presents conclusions summarizing this paper.

II. LUMPED-ELEMENT KINETIC INDUCTANCE DETECTOR DESIGN

The presented LEKIDs are based on series superconducting capacitor-inductor resonators coupled to a single transmission line, as it was first proposed in [11]. The inductor, which acts as the effective absorber, is designed to absorb the incident radiation by matching optically to free space. In this case, the strip grating, deposited on top of a high resistive silicon substrate, was designed to be matched at W band to the free-space impedance, when a backshort is placed at the rear part of the silicon wafer. The linearly polarized wave absorption of this design is characterized elsewhere and peaks around 78 GHz [17].

A new dual-polarization LEKID is designed to absorb and distinguish millimeter-wave radiation of two orthogonal linearly polarized waves simultaneously. For this purpose, two LEKIDs, based on the design explained before, can be stacked one on top of the other with perpendicular orientation as shown in Fig. 2. In this case, each of them will be dedicated to the detection of one of the two perpendicular polarizations of the incident waves.

Fig. 2. Cross section of a dual-polarization LEKID on silicon substrate.

A. Mm-wave Coupling Design

An equivalent circuit model of the dual-polarization LEKID is shown in Fig. 3.

Fig. 3. Equivalent circuit to model a dual-polarization LEKID. Short-circuit (SC) termination accounts for the presence of a backshort. The shunt admittances $Y_{\text{strip}}$ and $Y_{\text{strip2}}$ are the equivalent circuits of the strips.

The input admittance of the dual-polarization LEKID is

$$Y_{in} = Y_{\text{strip}} + \frac{1 + \tanh(y_f \ell_s)}{\eta_s} Y_1$$

(1)

where the admittance $Y_1$ is given by

$$Y_1 = Y_{\text{strip}} + \frac{1}{\eta_s \tanh(y_f \ell_1)}$$

(2)
and \( \gamma = \alpha_{Si} + j\beta_{Si} \) is the complex propagation constant, where the real part is the attenuation constant \( \alpha_{Si} \), due to dielectric losses, and the imaginary part \( \beta_{Si} \) is the phase constant in silicon. The shunt admittances \( Y_{strip1} \) and \( Y_{strip2} \) are the equivalent circuits of the strips that compose the inductors of the LEKIDs. Silicon substrates thicknesses are the equivalent circuits of the strips that compose the inductors.

The strips that compose the LEKID inductors have width \( 2a \), distance between strips \( g \) and a sheet resistance \( R_s \) (Ohm/sq). It is worth to note that even though the detector is operating at temperatures well below its critical temperature \( Tc \), an incident radiation with frequency \( (\nu) \) that provides energy higher than the superconducting gap \( (\Delta) \), i.e. \( h\nu > 2\Delta \), where \( h \) is the Planck constant, leads to a sheet resistance comparable to that in normal conducting state just above the superconducting critical temperature \[18\].

Metallic strips present a resistive and inductive impedance for an incident electric field parallel to them \[18\]-\[20\], which admittance is given by

\[
Y_{strip} = \frac{1}{R_s \left( \frac{g}{2a} + jX_L \right)} \tag{3}
\]

where the reactance \( X_L \) is

\[
X_L = \frac{\eta_0 g}{\lambda_0} \ln \left( \csc \left( \frac{\pi a}{g} \right) \right) \tag{4}
\]

If the incident electric field is orthogonal to the strips, the equivalent circuit is a capacitive admittance, according to \[21\], Sec. 5.18, eq.(1a).

A dual-polarization LEKID with two stacked and orthogonal strips is shown in Fig. 4. The admittance \( Y_{in} \) in (1) depends on incident wave polarization shown in Fig. 2. The admittance \( Y_{strip1} \) is inductive as (3) and \( Y_{strip2} \) is capacitive in case of a wave TEM \( V \). On the other hand, admittances \( Y_{strip1} \) and \( Y_{strip2} \) are capacitive and inductive as (3) respectively for TEM \( H \).

\begin{figure}[h]
\centering
\includegraphics[width=0.4\textwidth]{fig4.png}
\caption{Dual-polarization LEKID and two incident TEM waves with electric fields orthogonal between them. Each LEKID is composed of a meander line inductor and an interdigital capacitor.}
\end{figure}

The strip is designed to achieve efficient absorption, therefore, for an inductive strip admittance in (3), the geometry \((a, g)\) is calculated to fulfill \( \text{real} \left( Y_{strip1} \right) = 1/\eta_0 \) with metal resistance \( R_s (\text{Ohm/sq}) \) at cryogenic temperature. Then the grounded silicon substrate presents the imaginary admittance required to match the imaginary part of the admittance in (3) at the center frequency for a thickness \( \ell_1 \) given by

\[
\ell_1 = \frac{1}{\beta_{Si}} \left( \pi - \cot^{-1} \left( \frac{\eta_0 X_L}{\left( R_s \frac{g}{2a} \right)^2 + X_L^2} \right) \right) \tag{5}
\]

The second LEKID formed by rotating the same strips geometry \( 90^\circ \) with respect to the previous one on silicon substrate, with thickness a half wavelength at the center frequency, which allows matching simultaneously in both strips. This type of structure has narrow bandwidth behavior, since impedance matching is achieved with one single silicon dielectric layer \( \ell_1 \).

The initial geometry of the strips for the LEKIDs design has been obtained for a 35 nm thick Ti/Al film with \( R_s = 1.27 \) Ohm/sq measured just above \( Tc \) (see Fig. 5). A strip width \( 2a = 3 \mu m \) has been chosen to achieve the broadest bandwidth around 90 GHz, which together with a spacing \( g = 375 \mu m \) provides \( \text{real} \left( Y_{strip1} \right) = 1/\eta_0 \) in (3). Using those values, the inductive and capacitive admittances are calculated. Their reflection coefficients on air, from 65 to 110 GHz, are depicted in Fig. 6 (a).

\begin{figure}[h]
\centering
\includegraphics[width=0.4\textwidth]{fig5.png}
\caption{Sheet resistance versus temperature of a thin titanium/aluminum (Ti/Al) bilayer sample.}
\end{figure}

The thickness \( \ell_1 \approx 292 \mu m \) of the grounded silicon substrate is calculated using (5), and the silicon substrate thickness between LEKIDs is \( \ell_2 = \lambda_{Si}/2 \approx 483 \mu m \) with \( \lambda_{Si} = \lambda_0/\sqrt{\varepsilon_{rSi}} \), both calculated at 90 GHz. The obtained reflection coefficients, of the dual-polarization LEKID model in Fig. 2, for the two linear polarized incident waves (TEM \( V \), TEM \( H \)) and lossless silicon substrate, are plotted in Fig. 6(b) from 65 to 110 GHz, showing a perfect matching and maximum absorption at 90 GHz. The loss tangent of high resistive silicon in millimeter bands at 4.2 K is around \( 1.12 \times 10^{-6} \) \[22\], therefore it has been considered negligible for the design. Moreover, the obtained
impedance matching, for both incident waves, is not affected by
the strip with capacitive performance.

![Image](image1)

Fig. 6. Ti/Al film 35 nm thick strips with $R_s = 1.27$ Ohm/sq, $2a = 3 \ \mu m$ and $g = 375 \ \mu m$ in the frequency band 65 – 110 GHz: (a) Reflection coefficient of the strips on air, for incident linearly polarized waves parallel and orthogonal to the strips (b) Input reflection coefficient of the dual-polarization LEKID for both incident waves TEM.

With this initial design, the next step is a 3D electromagnetic simulation at W-band of the dual-polarized LEKID using the HFSS 3D electromagnetic simulator from ANSYS, as a single unit cell with Floquet ports with master and slave boundaries, to simulate an array as a planar-periodic structure. The final millimeter wave absorption, of the dual-polarization LEKID, for the two linearly polarized incident waves, is also simulated with the complete structure composed of the meandered strips and interdigital capacitors (see Fig. 4), and it is shown in [17]. The absorption efficiency in each component, using the Field Calculator available in the HFSS 3D electromagnetic simulator, is shown in Fig. 7 for both polarizations. This tool makes use of the Surface Loss function to estimate the energy absorbed in both inductors and capacitors in the LEKID [23]. The simulation results show a total absorption efficiency of 98% for each mode, showing the inductors the highest absorption efficiency with ~90%.

B. Low frequency Resonator Design

A typical LEKID superconducting resonator is coupled to the main microstrip readout line as it is shown in Fig. 1. The interdigital capacitor fixes its resonant frequency and allows frequency multiplexing. For readout, the resonator is coupled to a transmission line, selecting the coupling coefficient with a separation $s$ (see Fig. 1). Absorption of photons at millimeter wave results in a change of resonator frequency and quality factor, therefore an accurate design of the coupling is crucial for dynamic range optimization of the detector under millimeter wave radiation.

The connection of inductive and capacitive parts forms a series RLC resonant circuit. The resistive part of this equivalent circuit comes from conductor losses of both inductive and capacitive parts. As we are dealing with superconducting resonators, this resistive part can be extracted from the two-fluid model and tends to zero when the operating temperature is much lower than its critical temperature [11].

LEKID resonators are absorption-mode coupled [24], and the effect of an individual resonator on readout line transmission is only present in a narrow band around its resonant frequency. Fig. 8 shows an equivalent circuit with two coupled resonators in a transmission test schematic [25].

![Image](image2)

Fig. 7. Simulated absorption efficiency of each component, inductors (Ind.) and capacitors (Cap.), and total absorption $(1 - |S_{11}|^2)$ for the dual-polarization LEKID; (a) TEM, wave, with maximum absorption in the LEKID at the bottom. (b) TEM, wave with maximum absorption in the LEKID at the top.

Fig. 8. Equivalent circuit of two coupled resonators, in absorption mode, to the readout line. $M$ is the mutual inductance. $R_0$, $L_0$, $C_0$ represent the equivalent series resonant circuit of a LEKID.

Mutual inductance $M$ is the main coupling effect, which transforms each series-resonant circuit to a shunt-resonant circuit. In a narrow frequency band, taking into account only one resonator and considering other resonances well detuned, the equivalent circuit shown in Fig. 9 is useful to analyze LEKID resonator quality factors and coupling factor, and to obtain their relation to the $S_{21}$ scattering parameter. This last transmission parameter is in fact the tested one in a LEKIDs instrument.
Coupling factor $k$ is the relation between delivered power to external loads $P_E$ and internally delivered power $P_0$ inside the resonator, at resonant frequency, given by

$$k = \frac{P_E}{P_0} = \frac{R}{2Z_0} = \frac{Q_i}{Q_c}$$

(6)

where $Q_i$ is the internal quality factor and $Q_c$ the external quality factors [26]. For critical coupling, $k$ is 1 and the power delivered to the resonator is maximum.

LEKID coupling is analyzed as a pair of asymmetric coupled microstrip lines, having a total coupling length $\ell$ equal to the contribution of meander line sections close to the readout line: $\ell = 4\left(g + 2a\right)$, see Fig. 1. In a typical LEKID this coupling length is very short in terms of wavelength at the resonant frequency ($f_0$), and the equivalent circuit can be analyzed as a differential length ($dz$) of a coupled lines pair, as it is shown in Fig. 10(a). All elements ($C_1$, $C_2$, $C_{16}$, $L_1$, $L_2$ and $M$) have dimensions per unit of length. Readout microstrip line width ($W_1$) is typically chosen to have 50 Ohm characteristic impedance, whereas inductor strip width $W_2 = 2a$ is narrow in order to optimize the LEKID millimeter wave absorption, as well as to achieve an inductive value adequate for the resonant frequency, inside the selected radiofrequency band for the readout.

Performing a circuit analysis of the schematic shown in Fig. 10(b), using typical values for the microstrip structure, it is straightforward to realize that the most relevant element in the coupling is mutual inductance $M dz$. The equivalent circuit can be greatly simplified for design purposes, since all capacitive elements have a negligible effect. Using the equivalent circuits shown in Fig. 8 and Fig. 9, the mutual inductance required to obtain a specific coupling for a LEKID resonator is:

$$M_z = M \cdot \ell = \frac{1}{\omega_0^2} \sqrt{R R_0}$$

(7)

where $\omega_0$ is the resonant angular frequency $2\pi f_0$.

Coupling design can be simplified using symmetry properties of a two-port network [27]. Symmetry axis of a LEKID coupled resonator is depicted in Fig. 11. Applying odd mode condition, a short circuit in all points of the symmetry axis, the two-port network is converted to a one-port network, which is applied to the coupled lines network. Moreover, due to symmetry properties, the LEKID impedance $Z_{KID}$, the inductance elements in mutual inductance equivalent network and the coupling length are halved. The resultant equivalent networks are shown in Fig. 12(a) and Fig. 12(b). These two networks have identical $Y$ parameters, being their calculation the first step for the proposed design process. By analysis of mutual inductance circuit in Fig. 12(b) its $Y_{21}$ parameter is

$$Y_{21} = -j \frac{\omega_0}{\omega_0^2 \left(\frac{M_z}{Z} \right)^2 - \frac{L_1}{2} \frac{L_{2\ell}}{2}}$$

(8)

For a specific LEKID coupling design, the input data are: dielectric substrate parameters, the width of coupled lines, $W_1$ and $W_2 = 2a$ in Fig. 10(a), the coupling length $\ell$ and the resonant frequency $f_0$. Given the input data, and after selecting the desired coupling level: overcoupling, undercoupling or critical coupling, the mutual inductance $M_z$ is calculated according to (7). Inductances $L_{1\ell}$ and $L_{2\ell}$ in (8) are calculated from inductances per unit length $L_1$ and $L_2$ of isolated microstrip lines, for widths $W_1$ and $W_2$. The assumption of isolation of microstrip lines, for $L_1$ and $L_2$ calculation purposes, is a good approach, because coupling of LEKID resonators is very weak in general, and both lines are separated by a sufficiently large distance $s$.

The next step is to obtain the right separation $s$ between lines using an asymmetrical coupled microstrip lines electrical model [28]. This model is available in microwave circuit simulators, and through an optimization routine, or by manual tuning, the separation $s$ value must be varied to achieve the desired $Y_{21}$ parameter, of the two-port network in Fig. 12(c), with the value
calculated in (8) at the resonant frequency \( \omega_0 \). For an improved accuracy of the separation \( s \), a 2D electromagnetic simulator is used, taking as an initial value the obtained \( s \) value by electrical model simulation, to avoid the uncertainties of coupled lines electrical models for very weak couplings and short coupling lengths.

For the readout design of a LEKID, applying the present design method, Table I shows the obtained separation \( s \) for three different coupling factors regarding a LEKID with an expected internal quality factor \( Q_i = 200 \, 000 \) and resonant frequency \( f_0 = 500 \, \text{MHz} \). It has been considered a high resistive silicon substrate with thickness \( 285 \, \mu\text{m} \), coupling length \( \ell = 4 \cdot (g + 2a) = 3 \, \text{mm} \), and line widths \( W_1 = 220 \, \mu\text{m} \) for 50 Ohm readout line, \( W_2 = 2a = 3 \, \mu\text{m} \) for inductor strip. The \( S_{21} \) responses of a LEKID for the three coupling factors \( k \) with the parameters obtained in Table I are simulated and the results are shown in Fig. 13.

| Coupling | \( k \) | \( |S_{21}|_{\text{min}} \) (dB) | \( s \) (\( \mu\text{m} \)) | \( Q \) |
|----------|-------|-----------------|-----------------|------|
| over     | 3     | -12.0           | 323             | 50 000 |
| critical | 1     | -6.0            | 505             | 100 000 |
| under    | 0.25  | -1.94           | 787             | 160 000 |

Fig. 13. Simulated \( S_{21} \) response of a LEKID for three coupling factors \( k \) with the parameters shown in Table 1.

### III. KIDs Fabrication And Assembly

The devices were fabricated following the technological process detailed in [17]. In summary, by means of confocal sputtering, a Ti/Al bilayer was deposited on a high resistive silicon wafer \( 275 \, \mu\text{m} \) thick, with a \( 200 \, \text{nm} \) thick Al layer on the rear part of the wafer, to be used as ground plane and optical backshort. The physical dimensions for the inductor were calculated considering an initial approach of \( R_s \frac{L}{2a} \approx 377 \, \text{Ohm} \) in (3). The LEKIDs were designed for single polarization, and two prototypes were fabricated and measured at cryogenic and ambient temperatures with \( 2a = 3 \, \mu\text{m} \) and \( g = 440 \, \mu\text{m} \) for optical coupling. The bilayer was characterized at cryogenic temperatures, obtaining \( R_s = 1.27 \, \text{Ohm/sq} \), \( T_c = 782 \, \text{mK} \) and an estimation of the kinetic inductance \( L_k = 2.24 \, \text{pH/sq} \). Preliminary results for low frequency cryogenic characterization, and optical absorption for single polarization, were presented in [17].

In order to verify the performance of a dual-polarization prototype, several stacked wafers were measured at ambient temperature as a proof of concept. An array of \( 11 \times 11 \) single-polarization LEKIDs was fabricated on a \( 275 \, \mu\text{m} \) thick silicon wafer, while an identical second one was fabricated without a ground plane. Both wafers were stacked orthogonally with a \( 275 \, \mu\text{m} \) thick silicon wafer in between, to distinguish the polarization between two linear polarized incident waves. Fig. 14(a) shows a photograph of the stacked wafers placed in a test fixture, to characterize them at ambient temperature, whereas a photograph of the LEKID inductor is depicted in Fig. 14(b).

On the other hand, a new prototype with seven pixels was designed and fabricated modifying the coupling factor \( Q_c \) between pixels. Separation \( s \) and orientation of each LEKID, with respect to the single 50 Ohm microstrip transmission line, were modified from pixel to pixel in order to tune external coupling \( Q_e \). This device has been designed following the methodology explained in Section II.B, which is crucial for future designs, enabling the optimized critical coupling, i.e., \( Q_i \) to be close to \( Q_c \) under the desired optical load [16]. Fig. 15(a) and Fig. 15(b) show the design of two pixels, where their coupling has been modified by rotating \( 90^\circ \) the original pixel orientation, increasing the total coupling length \( \ell \). In the remaining pixels, coupling has been modified by tuning the separation \( s \) between lines. This prototype was mounted inside a light-tight package made from bulk aluminum, where aluminum wirebonds were used to connect microstrip lines to the readout chain as shown in Fig. 15(c).
Fig. 15. (a) and (b) show schemes of two of the pixels designed for coupling tuning. (c) Fabricated LEKIDs array mounted for cryogenic testing.

IV. EXPERIMENTAL SYSTEMS

This section describes two experimental test set-ups used for KIDs characterization: A W-band ambient temperature quasi-optical system used for millimeter wave absorption test and a cryogenic test system to measure LEKIDs resonant frequencies and quality factors.

A. Ambient Temperature Test System

The ambient temperature test system operates at W-band. It has been set up in order to characterize the LEKIDs absorption through a free-space measurement. The system consists of two horn antennas, two dielectric lenses that collimate the beam at the measurement plane, and a vector network analyzer [29], [30], [17]. The quasi-optical system has a 4f topology, with f'the focal length 75 mm of PTFE plano-convex lenses LAT075 (Thorlabs), shown in Fig. 16. The horn antennas (QSH-SL-75-110-F-20) provide around 20 dB gain. They are rectangular horns with dimensions \( a \times b = 14.8 \text{ mm} \times 11 \text{ mm} \) at the aperture, where the beam radii that maximize the coupling to the fundamental Gaussian mode are \( ax = 0.35a, ay = 0.5b \) [30]. The beam-waist radius at \( z = 0 \) is similar for both coordinates \( b_{\omega r} \approx 4.72 \text{ mm} \). The lenses diameter is \( D = 50 \text{ mm} \) and the beam radius at the antenna aperture is \( \omega \approx 5.37 \text{ mm} \), since \( D > 4\omega l \) the fractional power lost is lower than 3.10^{-4} for the fundamental Gaussian mode [30]. The calculated output beam waist radius, at the middle of the system, is \( b_{\omega r,\text{out}} \approx 16.87 \text{ mm} \), which defines the size of the beam spot for the absorption measurement of LEKID arrays at ambient temperature.

B. Cryogenic Test System

The experimental cryogenic system is the BlueFors Dilution Refrigerator LD-250 shown Fig. 17. This cryostat consists of a dilution refrigerator (DR) backed by a two-stage pulse tube cooler that provides 60 K and 4 K temperature stages, while the DR provides a 0.7 K stage and a variable 12 mK–1 K stage, where the detectors are mounted, Fig. 17(a). Electrical cryogenic characterization has been performed using the following read-out system, which connects the external warm electronics with the LEKID array at cryogenic temperatures.

1) Cryogenic Harness

The cryogenic harness scheme is shown in Fig. 17, which has been carefully chosen in order to minimize thermal loading between stages. Stainless steel inner and outer conductor coaxial cable is used from room temperature to the 4 K stage, being thermalized in the 60 K stage by a DC block. A 20 dB attenuator reduces the 300 K radiation and dissipates the power in the 4 K stage. Cupronickel (CuNi) coaxial is used from the 4 K stage to the 12 mK stage. Again, two DC Blocks and an extra 10 dB attenuator reduces the noise contribution from the warmer stages. Finally, a semirigid copper coaxial cable connects the last DC-Block with the LEKID package. Aluminum wirebonds are used to connect a microstrip board to the LEKID chip. On the return path, a copper coaxial cable connects the package with a DC Block at the 12 mK stage. Then, a superconducting NbTi coaxial cable carries the signal from this stage up to the 4 K stage where a SiGe Low Noise Amplifier (Caltech-CITLF3) amplifies the signal (15 dB gain). Finally, CuNi coaxial cable carries the signal to the output-port of the cryostat.

2) Readout system

In order to measure the transmission characteristic of the LEKID array, a readout system has been assembled in a chassis. It is composed of several coaxial modules such as SPDT (Single Pole Double Thru), LNA (Low Noise Amplifier), Attenuator, Power Splitter and Quadrature Modulator. Four SPDT RF
Switches enable switching between Vector Network Analyzer (VNA) connection and the I/Q demodulator, through a Measurement Mode Option and it also switches between two different DUTs (Device Under Test) through a Channel Selection Option. The VNA measurements work from 40 MHz to 2.6 GHz, and I/Q demodulation measurements from 500 MHz to 2.6 GHz.

For I/Q measurements, the LO signal is generated by a commercial PSG Analog Signal Generator, and is split using power splitter. A step attenuator is used as a variable attenuator to tune the readout signal power level to the required level for the LEKIDs. The signal coming from the LEKID cryostat output is amplified with a gain block consisting of three amplifiers (~43 dB gain), to reach the RF power required by an I/Q demodulator. The block diagram of the readout is in Fig. 18. The “IN” and OUT” ports are connected to the cryogenic set-up shown in Fig. 17.

Fig. 18. Readout block diagram for measuring the resonances of a LEKID array.

V. RESULTS

The measurements described in the following section refer to the fabricated prototypes that have been detailed in Section IV. Measurements at ambient and cryogenic temperatures have been carried out, and their results have been compared with the simulations.

A. Measurements at Ambient Temperature

The dual-polarization prototype, made up of 11 x 11 LEKIDs, was characterized at ambient temperature. This prototype was measured for two orthogonal and linearly polarized waves in the 65 to 110 GHz frequency band. Results are depicted in Fig. 19. These measurements were made using the quasioptical test-bench described in Section IV. A photograph of the sample under test is shown in Fig. 20. Simulations using the obtained parameters for the Ti/Al bilayers at ambient temperature ($R_s = 4$ Ohm/sq) are also included. The actual tested silicon substrate height ($h = 295$ µm) confirms a good fitting between simulated and experimental results for both polarizations, showing a maximum absorption around 75 GHz. The tolerance of the silicon wafer thickness is responsible for the frequency shift, and its real tested value has been updated in the simulations.

Moreover, these measurement results have been compared with the simulated absorption of all LEKID components, in order to obtain the absorption efficiency. Fig. 19(a) and Fig. 19(b) show a comparison between test results and simulations of each component, which confirms that incident power is mainly absorbed in the inductor, as expected. The inductor absorption efficiency is reduced due to the dielectric loss, which has a negligible effect at cryogenic temperature.

Fig. 19. Absorption measurements and simulations of 11 x 11 LEKIDs at W-band for two orthogonal polarizations at ambient temperature; (a) TEM$_V$ wave, with maximum absorption in the LEKID at the bottom. (b) TEM$_H$ wave with maximum absorption in the LEKID at the top.

Fig. 20. Dual-polarization LEKID array under absorption test at W-band.
B. Cryogenic Measurements.

Cryogenic characterization of the 7 pixels array was performed using the experimental set-up explained in Section IV.B. A Vector Network Analyzer is employed for measuring the $S_{21}$ parameter across the array. Fig. 21 shows the $S_{21}$ transmission amplitude through the array, where each minimum corresponds to a pixel. Due to the low pixel packing density in this prototype, negligible crosstalk between resonators is expected [31]. The resonant frequency, loaded quality factor ($Q_l$), external quality factor ($Q_c$) and internal quality factor ($Q_i$) was obtained following the procedure detailed in [32].

![Fig. 21. Transmission sweep ($S_{21}$) measured for the fabricated Ti/Al LEKID array at T=12 mK and dark environment.](image)

The EM simulation software Sonnet was used to confirm the estimated kinetic inductance ($L_k$) from electrical characterization [17]. Fig. 22 shows the simulated resonant frequency of the lowest resonance as a function of the kinetic inductance. As can be seen, the experimental resonant frequency corresponds to $L_k = 2.2 \, \text{pH/sq}$, very close to the estimated value. These simulations estimate the kinetic fraction by comparing the experimental resonance and the simulated one, obtaining $\alpha = 0.38$.

![Fig. 22. Simulated resonant frequency as a function of the kinetic inductance. Vertical blue line indicates the measured experimental resonant frequency at 12 mK. Crossover between this line and the simulated values (red) shows the obtained kinetic inductance 2.2 pH/sq.](image)

The quality factors obtained for all of the pixels are shown in Fig. 23. The internal quality factor, around $10^6$ for all the pixels, indicates the excellent quality of the deposited film [33]. The external quality factor, as can be seen, has been tuned within an order of magnitude, following the procedure explained in Section II.B. The lowest value, $3.5 \cdot 10^4$ corresponds to the pixel shown in Fig. 15(a), whereas the highest value, $5 \cdot 10^6$, corresponds to the pixel shown in Fig. 15(b). The obtained external quality factors agree reasonably well with the values from the design methodology. Differences between measurement and simulation are related to small tolerances in the nanofabrication process.

![Fig. 23. Quality factor distribution for the 7 pixels fabricated, obtained at T=12 mK and dark environment. The obtained $Q_c$ values have been ordered from lowest to highest for clarity.](image)

Under these dark conditions, the external quality factor limits the loaded quality factor (overcoupling, $Q_l > Q_c$). However, under the optical load, the internal quality factor is expected to diminish, leading the LEKIDs to a critical optical coupling ($Q_c = Q_i$) which is desirable for maximizing the response [16]. For simulating this effect, a temperature sweep has been performed as shown in Fig. 24(a). Fig. 24(b) shows the loaded, internal and external quality factors as a function of bath temperature. As can be seen, $Q_c$ is reduced as temperature is increased due to thermally excited quasiparticles. This effect can be compared with the response of the LEKIDs upon optical illumination [34]. As can be seen, the LEKID goes from an overcoupled to undercoupled regime, showing the importance of a proper microwave design for reaching the critical coupling ($Q_c = Q_i$). Future experiments will be performed under optical illumination in order to choose the optimum external coupling.
observed. A crossover from overcoupling to undercoupling can be observed. (b) Loaded, internal and external quality factors as a function of bath temperature. A crossover from overcoupling to undercoupling can be observed.

VI. CONCLUSION

This work has demonstrated dual-polarization LEKID array absorption at W-band at ambient temperature. A superconducting Ti/Al bilayer is used to push down the critical temperature (from 1.2 K of pure Al to 782 mK), and therefore push down the low frequency limit imposed by the superconducting gap to this frequency band. The used topology allows us to detect simultaneously two orthogonal linearly polarized waves in one single pixel. The described design methodology is applicable to other millimeter and submillimeter wave bands.

On the other hand, a prototype with seven pixels was fabricated adjusting the external coupling factor, following a proposed method, in order to achieve critical coupling under the desired optical load. The results obtained exhibit a high internal quality factor. The range of the external quality factors applied has experimentally demonstrated a critical coupling when the bath temperature is increased to emulate an optical load. The dual-polarization LEKID design presented and its initial test results show a very promising technology for future polarimetry experiments.

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