Learn2Hop: Learned Optimization on Rough Landscapes
With Applications to Atomic Structural Optimization

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Abstract
Optimization of non-convex loss surfaces containing many local minima remains a critical problem in a variety of domains, including operations research, informatics, and material design. Yet, current techniques either require extremely high iteration counts or a large number of random restarts for good performance. In this work, we propose adapting recent developments in meta-learning to these many-minima problems by learning the optimization algorithm for various loss landscapes. We focus on problems from atomic structural optimization—finding low energy configurations of many-atom systems—including widely studied models such as bimetallic clusters and disordered silicon. We find that our optimizer learns a ‘hopping’ behavior which enables efficient exploration and improves the rate of low energy minima discovery. Finally, our learned optimizers show promising generalization with efficiency gains on never before seen tasks (e.g. new elements or compositions). Code will be made available shortly.

1. Introduction
Efficient global optimization remains a problem of general research interest, with applications to a range of fields including operations design (Ryoo & Sahinidis, 1995), network analysis (Abebe & Solomatine, 1998), and bioinformatics (Liwo et al., 1999). Within the fields of chemical physics and material design, efficient global optimization is particularly important for finding low potential energy configurations of isolated groups of atoms (clusters) and periodic systems (crystals); identifying low energy minima in these cases can yield new stable structures to be experimentally produced and tested for a wide variety of industrial or scientific applications (Wales & Doye, 1997; Flikkema & Bromley, 2004). However, even simple examples with a few atoms have quite complex minima structures. Numeric approximations suggest that systems of only 147 atoms could have $10^{60}$ distinct minima (Tsai & Jordan, 1993). Global optimization problems can also be quite difficult when high loss barriers exist between local minima, as depicted in Figure 1. Despite being NP-hard in the worst case, significant work has been put into developing optimization techniques for these structure prediction tasks. Nonetheless, classical approaches to this problem continue to face a number of drawbacks including requirements of: a significant number of steps (Wales & Doye, 1997; Pickard & Needs, 2011), carefully selected hand-crafted features, or sensitive dependence on learning rate schedules (Bitezek et al., 2006).

In this work, we propose adopting a new class of strategies to these global optimization problems: learned optimization (Bengio et al., 1992; Andrychowicz et al., 2016; Metz et al., 2018). Here, hand-designed update equations are replaced with a learned function parameterized by a neural network and trained via meta-optimization. While this strategy has shown promise for training neural networks (Metz et al., 2020) where falling into bad local minima is not a concern (Choromanska et al., 2015; Luo et al., 2018), current techniques fail to prioritize global minimum discovery or have not been tested on rough loss landscapes with many unconnected local minima.

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substantial improvement over classical algorithms for these sorts of global optimization problems. To this end, we present several modifications of learned optimizers required to effectively train models which can find low energy states of many-minima loss surfaces. Using several canonical problems in atomic structural optimization, we demonstrate that the learned optimizers outperform their classical counterparts when trained and tested on similar systems and—more surprisingly—are able to generalize to unseen systems. The specific contributions of this paper are as follows:

1. Novel parameterizations of learned optimizers prioritize global minimum discovery (Section 3) and yield improvements on benchmark tasks consisting of single atom types (Section 4).
2. Analysis of learned optimizer behavior showcases an automatically-learned ‘hopping’ behavior which enables efficient exploration of minima (Section 5).
3. Results for bimetallic systems show that our learned optimizers can generalize beyond the examples seen during training, yielding gains in efficiency and performance over commonplace optimization techniques such as Basin Hopping (Section 6).

2. Background / Related Work

2.1. Atomic Empirical Potentials

Atomic structure optimization often requires finding the lowest energy configuration of a system (Oganov et al., 2019). However, accurate calculation of energies is expensive, often requiring quantum mechanical simulations such as DFT (Hohenberg & Kohn, 1964). In this work, we instead use approximations of the potential energy known as empirical potentials, that are not only simple and efficient to calculate but also have minima that correlate to those found by more accurate calculations (Tran & Johnston, 2011). The empirical potentials studied in this paper are as follows:

**Lennard-Jones Clusters** are often used in the modelling of spherically-symmetric particles in free-space such as noble gasses or methane and are the archetypal model for a simple-to-compute potential (Jones & Chapman, 1924; Doye et al., 1999). The total energy of the system is defined only by pairwise distances, denoted \( d_{ij} \):

\[
\sum_{i} \sum_{j > i} \epsilon \left[ \left( \frac{d_0}{d_{ij}} \right)^{12} - 2 \left( \frac{d_0}{d_{ij}} \right)^{6} \right]
\]  

where \( \epsilon \) is the minimum two-particle energy and \( d_0 \) is the distance where this occurs. Despite the apparent simplicity, the minima structures of these systems are complex and vary significantly based on the number of atoms (Doye et al., 1999; Wales & Doye, 1997). For example, the 13 and 19 atom systems display concentrated “funnels”, where many local minima and the global minimum are connected via low energy paths. In contrast, the 38 and 75 atom counts display complex “double-funnelled” landscapes, where there are two distinct sets of minima that are dynamically inaccessible due to a high energy barrier between the two.

**Gupta Clusters** are similar to the Lennard-Jones model in approximating the energy of sets of atoms in free-space, yet they are significantly more complex due to the inclusion of a second-moment approximation of the tight-binding Hamiltonian (Gupta, 1981; Michaelian et al., 1999; Sutton, 1993; Rosato et al., 1989). In this paper, we focus on both single element and bimetallic forms using the elements Ag, Au, Pd, and Pt. Energy equations and the constant values for all systems are provided in Appendix A.2.

**Stillinger-Weber (SW)** potentials (Stillinger & Weber, 1985) provide more accurate estimations for energies of semiconductors. This empirical potential introduces a three-dimensional potential that is more complex and accurate than the Lennard-Jones and Gupta potentials.

Diagrams depicting the local minima structures and the lowest energy paths between minima for particularly interesting cluster sizes can be viewed at [http://doye.chem.ox.ac.uk/research/forest/lj.html](http://doye.chem.ox.ac.uk/research/forest/lj.html)
body angular term between atoms, making the corresponding loss landscape significantly more difficult to optimize. In this paper, the SW potential is used to model silicon crystals. This benchmark is distinct from the others in the use of periodic boundary conditions, so that atomic structures are tiled in space. The associated energy equation and parameters are provided in Appendix A.3.

2.2. Optimization Methods from Structure Prediction

Early approaches to structure prediction problems simply initialized the particle positions at hand-crafted, physically-motivated structures, before applying gradient descent (Hoare & Pal, 1971; Farges et al., 1985; Doye et al., 1995). This technique proved effective for simple cluster systems such as Lennard-Jones but faced difficulty scaling to more complex potentials (such as those with angular dynamics). Classic optimization approaches such as Basin Hopping (Wales & Doye, 1997) and Simulated Annealing (Kirkpatrick et al., 1983; Biswas & Hamann, 1986) resulted in significant improvements and helped discover the minima for a variety of structures. However, these techniques end up requiring high step counts and may only find the global minimum in the limit of infinite optimization steps.

Modern molecular dynamics systems use a variety of techniques for optimization. Quasi-Newton techniques such as BFGS and damped Beeman dynamics (Beeman, 1976) are popular within libraries such as QuantumEspresso (Gianozzi et al., 2009). Alternate strategies include Fast Inertial Relaxation Engine—referred to as FIRE (Bitzek et al., 2006)—which adaptively modifies the velocity over the course of training and Ab Initio Random Structure Search (Pickard & Needs, 2006; 2011; Zilka et al., 2017). However, these techniques often rely on heuristics or require a large number of restarts before reaching the global minimum.

While this work only uses traditional empirical potentials, machine learning has also been used to create empirical potentials, such as those utilizing graph convolutions (Gilmer et al., 2017; Schütt et al., 2017; Cheon et al., 2020). These models are becoming a popular option for speeding up optimization. However, we note that the approach presented in this paper is complementary; the two could be combined so that both the potential and optimizer are learned.

2.3. Learned Optimization

Learned optimization (Bengio et al., 1992; Andrychowicz et al., 2016; Wichrowska et al., 2017; Lv et al., 2017; Metz et al., 2018; 2019b; Gu et al., 2019; Metz et al., 2020), has recently become a popular meta-optimization task, where updates are a function of the gradients, parameterized by a neural network. In the traditional setup depicted in Figure 2, training a learned optimizer consists of an inner-loop of optimization problems which are used to compute meta-updates to the learned optimizer parameters, referred to as the outer-loop (Wichrowska et al., 2017; Metz et al., 2018).

In our case, the inner-loop consists of instantiations of atomic structure optimization problems, including a random initialization for atoms and a corresponding empirical potential to minimize. At each step in the inner-loop of meta-training, atomic forces are computed, featurized, and then input to the learned optimizer which computes updates to the particles. These steps are then repeated, which is often referred to as an inner-trajectory or unroll.

For each inner-loop, a meta-loss is defined based on the optimization trajectory, commonly the average loss over the trajectory in prior work. If the unrolls were short, meta-training could be performed by gradient descent (Andrychowicz et al., 2016; Wichrowska et al., 2017). However, due to memory requirements and often ill-conditioned outer-loss surfaces (Metz et al., 2019a), meta-gradients are instead approximated via Evolutionary Strategies (ES) using antithetic samples (Williams, 1992; Salimans et al., 2017; Metz et al., 2019a). A central controller collects batches of meta-gradient estimates and updates the learned optimizer parameters, typically using Adam (Kingma & Ba, 2014).

A variety of architectures have been proposed for learned optimizers. Early work utilized RNNs in order to provide the network a state that can be automatically updated throughout the course of training (Andrychowicz et al., 2016). These models were quickly developed into optimizers that scale (Wichrowska et al., 2017), and are compute-efficient (Metz et al., 2018). Most closely related to our work is that of (Metz et al., 2019a; 2018) which is novel in its parameterization of the state and input features of the learned optimizer. Instead of providing an explicit memory (e.g. in a GRU), the learned optimizer is simplified to an MLP that is applied per parameter and is provided relevant features, such as the first and second moment estimates for the gradients. See Table 1 for all features used in the MLPs.

3. Modifications for Rough Landscapes

Adapting these learned optimizers to many-minima landscapes requires modifications to both the training and model itself to improve global optimization. For example, instead of the average loss of the optimization trajectory, the learned optimizer for atomic structure optimization only uses the final step loss. This training strategy prioritizes global minimum discovery at the expense of greater variance of the gradient with respect to learned optimizer weights. Additional modifications are detailed in the following sections and training details are in Appendix B.
Table 1. Features inputted into the learned optimizers. MLPOpt refers to the model by Metz et al. (2019a).

| FEATURES        | DESCRIPTION IN MLPOPT |
|-----------------|------------------------|
| Gradients       | Gradients              |
| Positions       | Particle positions     |
| Decays          | EMA of 1st and 2nd moments |
| Adam-Like       | Moment correction      |
| Singular        | Number of particles    |
| Species         | Species identity       |
| Step            | Training step sine features |
| Radial          | Radial symmetry features |

3.1. Features

We follow prior work (Metz et al., 2019a) in parameterizing the learned optimizer as an MLP. The input features are often inspired by popular optimization techniques and include estimates of the first and second moments to mimic Adam (Kingma & Ba, 2014). Table 1 describes all inputted features that are adopted from “MLPOpt”, the learned optimizer described in Metz et al. (2019a).

Novel to our learned optimizers is the inclusion of Behler-Parrinello radial symmetry features (Behler & Parrinello, 2007; Artrith et al., 2013; Cubuk et al., 2015). Traditional learned optimizers update each parameter independently, yet in the case of atomic structure problems, particle behavior should depend on interactions with nearest neighbors. Radial symmetry functions provide these sorts of two-body interactions for a central atom by allowing updates to be defined by local neighborhoods. Simply put, these features \( \phi \) are computed using a Gaussian kernel and summing over all neighbors of a central atom. Smooth cutoffs \( \Gamma \) are applied using the formulation by Behler & Parrinello (2007):

\[
\phi_i = \sum_{j \neq i} \exp \left( -\eta d_{ij}^2 \right) \Gamma(d_{ij})
\]

\[
\Gamma(d_{ij}) = \begin{cases} 
0 & \text{if } d_{ij} > c \\
0.5 \left( \cos \left( \pi \cdot \frac{d_{ij}/c}{1} \right) + 1 \right) & \text{otherwise}
\end{cases}
\]

where \( c \) is a pre-defined cutoff set to 2.5 angstroms and \( \eta \) is a hyperparameter controlling the scale. \( \eta \) is set to one of \{0.0009, 0.01, 0.02, 0.035, 0.06, 0.1, 0.2, 0.4\}, yielding 8 radial features per atom type.

These features are then parameterized into a log-magnitude and direction representation (Andrychowicz et al., 2016):

\[
\text{features} = \begin{cases} 
\left( \log|x|/p, \text{sign}(x) \right) & \text{if } x > \exp(-p) \\
\left( -1, x \exp(p) \right) & \text{if } x \leq \exp(-p)
\end{cases}
\]

where \( p = 10 \) is the default hyperparameter. These features are then input to the learned optimizer, a small 2-layer dense neural network (with a hidden size of 32), that acts component-wise. The network outputs magnitude \( m \) and unnormalized direction \( d \) per component, converted to the final update via:

\[
\alpha \cdot d \cdot \text{sigmoid}(\beta \cdot m + \gamma)
\]

where \( \alpha, \beta, \gamma \) are scalars learned via meta-optimization.

3.2. Meta-Training Stability

The rough loss landscapes discussed in this paper present two significant challenges with regards to meta-optimization: high curvature and infrequent training signal.

High curvature is an intrinsic problem to atomic structures. For example, with the Lennard-Jones potentials, the energy increases at a rate of \( d_{ij}^{-12} \) as \( d_{ij} \to 0 \) for all \( i, j \). When the optimizer happens to bring two particles too close together, energy (loss) spikes can yield meta-gradients that destabilize learned optimizer training. Traditional strategies such as gradient norm clipping (Pascual et al., 2013) were found to be ineffective in preventing divergence of the meta-optimization objective.

Infrequent and noisy training signals are also problematic as learned optimizers can find simple, stable optimization strategies such as gradient descent early in training. Most perturbations to gradient-descent methods will be noise and increase the final loss. The meta-optimization model must be sensitive enough to learn from the infrequent signal occurring when few individual instantiations of a learned optimizer find better minima structure, rather than being pushed back to descent-like methods due to noise.

As mentioned in the background, many learned optimizers are trained with antithetic ES sampling (Salimans et al., 2017; Metz et al., 2019a) where meta-gradients are estimated via perturbations of the parameters:

\[
\nabla_{\theta} L \approx \mathbb{E}_{\epsilon \sim \mathcal{N}(0, I\sigma^2)} \left[ \frac{L(\theta + \epsilon) - L(\theta - \epsilon)}{2\sigma^2} \right] \epsilon
\]

where \( L \) is the loss, \( \theta \) the learned optimizer parameters, and \( \sigma \) is the perturbation scale, set to 0.1. This strategy is particularly vulnerable to the optimization difficulties, as either direction of the parameter perturbation may lead to exploding gradients. Also, the variance of the estimator makes it more difficult to learn from the sparse rewards when optimizers find better local minima. To overcome these issues, we present two modifications to the meta-update which enable stable meta-training:

\( ^3 \text{Note, this output parameterization contrasts from Metz et al. (2019a), but experimental evidence showed that the traditional exponential update leads to divergent optimization trajectories.} \)
Meta-loss Clipping (ESMC)

In order to prioritize signal from perturbations that find better minima and improve the meta-loss, we propose clipping the loss functions in the meta-gradient computation at the value found by the unperturbed parameters.

$$E_\epsilon \left[ \min \left[ L(\theta), L(\theta + \epsilon) \right] \right] - \min \left[ L(\theta), L(\theta - \epsilon) \right] \right] = \frac{L(\theta + \epsilon) - L(\theta - \epsilon)}{2\sigma^2}$$

where again $\epsilon \sim \mathcal{N}(0, I\sigma^2)$.

Intuitively, this biases against directions of high curvature in meta-optimization and empirically showed improved results. This strategy has the added benefit of heavily clipping the gradients of examples where loss spikes when atoms become too close, at the cost of an additional meta-loss calculation for $L(\theta)$.

Genetic Algorithms (GA)

Instead of relying on approximate meta-gradients, a simpler strategy perhaps is to adopt the perturbed parameters when they improve the meta-loss on a batch of random examples (Holland, 1992; Goldberg & Holland, 1988). To match the number of estimators of the meta-gradient used in ESMC, we use a population of size 80. At the end of each outer loop, the best performing parameters $\theta$ are kept and used for creating the next population by drawing from $\mathcal{N}(\theta, I\sigma^2)$ where $\sigma = 0.1$. By default, $\theta$ is kept constant when all samples perform worse than the baseline.

Comparison of Methods

A comparison of these strategies on a simplified learned optimizer setup is shown in Figure 3. The genetic-algorithm approach shows improvement early in meta-training which steadily converges to an optimizer where almost all initialization find the global minimum. In contrast, both ES and ESMC show a distinct transition in behavior around steps 300-400, which demarcates a transition from simple-to-learn descent behavior to more complex global minima discovery. The ESMC method is able to retain this behavior throughout meta-optimization, whereas traditional ES appears unstable and forgets. Overall, both learned optimizer modifications show significant improvements in convergence speed and stability when compared to vanilla ES. Details for this experiment can be found in Appendix B.

3.3. Additional Details

In order to provide consistent scaling when averaging meta-losses across tasks, we divide energies by the best minimum found from applying Adam to 150 random initializations. This normalizes all losses to so that $-1$ is the best minimum found by Adam, ensuring that optimizers trained on multiple particle counts are not biased towards larger systems where the energy scales are lower. For each inner-loop, we apply 50000 optimization steps before computing meta-gradients. Batched training occurs with 80 random initializations of atomic structure problems. Once meta-gradients are averaged, a central controller meta-updates the parameters of the learned optimizer via Adam with a learning rate of $10^{-2}$ (which decays exponentially by 0.98 every 10 steps). This repeats for a total of 1000 meta-updates.

Finally, as the learned optimizer training does not guarantee a local minimum is found at the end of an optimization trajectory, we add 1000 steps of GD at learning rate of 0.001. We evaluate all strategies using 150 random initializations and report the mean and minimum energies found.

3.4. Implementation Details

The aforementioned potentials are coded using JAX-MD (Schoenholz & Cubuk, 2019). The learned optimizers are built in JAX (Bradbury et al., 2018) to take advantage of automatic differentiation and vectorization of the optimization simulation. The associated training and evaluation utilized V100 GPUs. For distributed training, the controller batches computation on up to 8 GPUs.

4. Experimental Results - Single Atom

We start with a simplified set of potentials comprised of a single type of atom. For Lennard-Jones, we present results when a model is trained on a diverse set of atom counts, specifically $\{13, 19, 31, 38, 55, 75\}$, which helps stabilize learned optimizer training and improve generalization beyond the training set. In the results, starting with Figure 4, the learned optimizer shows significant performance gains when compared to the benchmark optimization algorithms of Adam and FIRE. This improvement not only takes the form of better minima but also better average energy per
Figure 4. Comparison of the learned optimizers and baseline methods on Lennard-Jones clusters. The learned optimizers are trained on a subset of the atom counts (demarcated by dashed lines) and evaluated via 150 random initializations. When compared to Adam and FIRE, the learned optimizers generically improve both the average energy per initialization (top) and best minima found (bottom) on atom counts unseen during training.

Figure 5 analyzes the distribution of minima in greater detail for two canonical tasks: the 13 and 75 atom Lennard-Jones systems. The loss surface of the former is best described as a “funnel” and even traditional algorithms can find the global minimum in about 20/150 random initializations. On the other hand, the 75 atom Lennard-Jones system has a glassy, optimization landscape, where high energy barriers exist between local minima (Wales & Doye, 1997) and the global minimum is difficult to find.

Interestingly, we first find that the baselines of Adam and FIRE yield similar performance per task after extensive hyper-parameter tuning. Nevertheless, both of our learned optimizer models show significant progress. With 13 atoms, the learned optimizers drastically increase the rate of global minimum discovery from 20/150 to above 140/150. With 75 atoms, the learned optimizers shift the distribution of minima found, finding minima within 1 eV of the global minimum compared to the 7 eV for Adam and FIRE.

Similar results were obtained when learned optimizers were extended to the Gupta or SW potentials, when modeling 55-atom gold clusters and 64-atom silicon crystals respectively. Table 2 shows that the learned optimizers routinely surpass Adam and FIRE baselines and outperform Basin Hopping in a step-matched comparison.

5. Behavioral Analysis

As the update function is parameterized by a neural network, it is unclear how learned optimizers improve atomic structural prediction. To explore the behavior, we provide three analyses showcasing an emergent ‘hopping’ behavior and what features are critical for learned optimizer performance.

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6 Step-matched refers to an equivalent number of inner optimization steps. For details, please see the Appendix C.2

7 In this comparison, Basin Hopping is limited to 50000 steps.
5.1. Loss Trajectories

In Figure 6 (top), we show loss trajectories when the learned optimizer is applied to the Lennard-Jones task with 13 atoms. Interestingly, the behavior of the loss function is not monotonic. While the model does rapidly descend into individual basins, many of these models display spikes in loss or ‘hopping’ behavior where the model transitions between basins of different local minima at an erratic interval. More over, the optimizers have discovered characteristics that determine when to leave their basin. Figure 6 (bottom) explores these trajectories in greater details, by filtering the ‘lucky’ initializations that lead to the correct global minimum via Adam only. In cases where the parameters start in the correct basin, the learned optimizers performs better, acting like traditional Adam. For worse random starts, the learned optimizer will descend and then ‘hop’ between basins.

5.2. Behavior at Minima

This ‘hopping’ behavior appears to be key to the learned optimizer performance. Inspired by Maheswaranathan et al. (2021), we fix the positions at the global minimum and compute the learned optimizer updates over an entire optimization trajectory. This strategy removes the influence of gradients in the learned optimizer (as they are zero) and helps visualize behavior as a function of step number. In Figure 6 (middle), repeatedly applying the learned optimizer and measuring the update magnitude displays these ‘hops,’ indicating that the model emphasizes on exploration and hopping between basins midway through these optimization trajectories (despite not receiving gradient signal to do so).

5.3. Feature Importance

Table 3 provides an ablation study to explain what input features to the learned optimizer are most important. To clarify the difference in performance, results are presented for a learned optimizer trained only on the 75 atom Lennard-Jones system (similar results were found for other clusters and crystals). To account for training instability, each result is the median over 10 shortened runs of 650 meta-updates.

We first see slight improvements coming from the addition of exponential moving averages of the gradient, similar to the benefits of momentum in stochastic optimization. However, what is most critical to model performance is the training step sine features and the radial symmetry features. The sine features encode the optimization step via sine waves of various timescales \{1, 3, 10, 30, 100, 300, 1000, 3000, 10000\}, and we hypothesize that these are helpful as they provides signal for when the learned optimizer should perform exploration and exploitation. Otherwise, the models tend to learn monotonic behavior that is similar in spirit to the Adam solutions. Finally, the newly introduced radial symmetry features also provide significant improvement, suggesting that per-position optimization is sub-optimal and rather information about particle neighbors (other than just gradients) is informative for optimization.

Overall, the behavioral analyses show that ‘hopping’ be-

| POTENTIAL      | EVALUATION # ATOMS | GM     | METRIC | STEP-MATCHED BASELINES | LEARNED OPTIMIZER |
|----------------|--------------------|--------|--------|------------------------|-------------------|
| Lennard-Jones  | 13                 | -44.33 |        | -44.33                 | -44.33            |
|                | 75                 | -397.49|        | -40.58                 | -40.45            |
|                |                    |        | Min    | -390.34                | -389.12           |
|                |                    |        | Mean   | -380.52                | -380.23           |
| Gupta Gold     | 55                 | -181.89|        | -180.94                | -181.75           |
|                |                    |        | Min    | -179.94                | -181.38           |
|                |                    |        | Mean   | -175.85                | -176.38           |
| SW Silicon     | 64                 | -277.22|        | -60.08                 | -261.44           |
|                |                    |        | Min    | -256.83                | -257.01           |
|                |                    |        | Mean   | -258.46                | -258.14           |

| OPTIMIZERS     | MINIMUM ENERGY (eV) |
|----------------|----------------------|
| Baseline       |                     |
| Adam           | -390.3               |
| LEARNED OPTIMIZER |                 |
| ESMC (OURS)    | -390.3               |
| (1) GRADIENTS  | +0.2                 |
| (2) POSITIONS  | +0.2                 |
| (3) DECAYS     | -2.0                 |
| (4) ADAM-LIKE  | -1.0                 |
| (5) SINGULAR   | -1.0                 |
| (6) SPECIES    | -0.5                 |
| (7) SINE       | -2.3                 |
| (8) RADIAL     | -3.4                 |

Table 2. Learned optimizers show improvement across all tested potential types. For each model and evaluation, the average and minimum are computed across 150 random initializations. All energies are reported in units of eV, and GM denotes the global minimum energy.
Behavioral analyses of learned optimizers find that examples ‘hop’ between basins rather than descent behavior. This is seen in the behavior of individual trajectories (top), where each trajectory has an opacity of 0.02 (so darker regions correspond a greater number of examples). These ‘hops’ are further supported by the spiking update magnitudes when fixed to a local minimum, suggesting that learned optimizers prioritize exploration of various basins (middle). For both of these diagrams, we zoom-in on a single trajectory, showing how the ‘hops’ arise from large updates, followed by periods of descent. We also find that this hopping behavior corrects unlucky, random initializations that would not find the global minimum via Adam (bottom).

Behavior is critical to the performance of learned optimizers; however, given that an equivalent number of traditional Basin Hopping steps yields worse performance suggests more complex behavior also occurs.

### 6. Experimental Results - Bimetallic Clusters

Having found that learned optimizers perform well in the case of single atom systems, we introduce additional complexity and explore generalization performance of the learned optimizers using bimetallic clusters. These systems are particularly interesting as purely gradient-based optimization methods such as Adam or FIRE fail, unable to pass the large energy barriers between local minima. These potentials also allow for exploration of whether the learned behavior can transfer, a promising sign for the usage of these

models in material design or crystal discovery.

For the bimetallic clusters, we focus on the Gupta potential, whose parameters are modified to correspond to specific pairwise interactions (Gupta, 1981). The constant values used can be found in Appendix A.2. First, the results of training learned optimizers on bimetallic clusters comprising of Silver (Ag) and Gold (Au) are shown in Figure 7 (top). Fixing the total number of atoms at 38, we train the learned optimizer on $\text{Ag}_{38-m}\text{Au}_m$ for $m \in \{3, 7, 15, 22, 30, 38\}$ and test on all values of $m$. We present the convex hull of the formation energies of the clusters, as in equilibrium when excess silver and gold particles are present, only these clusters will be stable. The graphs show the robust empirical performance of learned optimizers, significantly outperforming Adam, FIRE, and a step-matched Basin Hopping benchmark. Only after 10x evaluation steps does Basin Hopping compete with the performance of the learned optimizer. This result indicates that the learned optimizers generalize as few of the AgAu clusters were used in training.

In the context of material design and crystal discovery, another core question is whether the learned optimizers will
generalize beyond the set of atoms used for training. In Figure 7 (bottom) we show the results from both the AgAu model described above and a second model trained on AgAu, AgPt, and PdAu. For both, we test on clusters of PdPt which is not in the training set of either model. The learned optimizers show successful transfer performance, exceeding the step-matched Basin Hopping results. Only after 10x the number of evaluation steps can Basin Hopping compete with the learned optimizers (even after tuning, see Appendix C.2. While increasing the diversity of training tasks does improve generalization performance, both optimizers show an ability to transfer to unseen elements or combinations, a promising sign for this strategy of learning to optimize.\(^8\)

7. Conclusion

With current optimization techniques in material design and chemical physics requiring hand-crafted features or significant evaluation time, this paper explores the idea of global minimum discovery using learned optimizers. Although novel adaptations are required from the current state-of-the-art learned optimizers, we show that the resulting models can beat current baseline optimization techniques such as Adam and FIRE, not only in terms of minima discovery but also in terms of average energy per initialization. Better yet, these learned optimizers show signs of transference between potentials of similar varieties (e.g., clusters parameterized by the Lennard-Jones and Gupta potentials), even on never before seen elements or combinations. Although a single optimizer for all task remains a goal for future work, learned optimizers show promise in automatically finding minima in complex optimization landscapes. We hope that the resulting models can aid in the design of new materials (e.g., for addressing energy challenges).

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A. Empirical Potentials

Empirical potentials are approximations of the potential energy surface for physical systems. In contrast to Density-Functional Theory which computes energies from first principles (Hohenberg & Kohn, 1964), empirical potentials are generic functional forms whose parameters are fitted to correspond to individual systems and which are designed to be efficient to compute. In our paper, these empirical potentials enable the training of learned optimizers, as we can easily batch computation and use automatic differentiation techniques to quickly calculate gradients/forces.

Moreover, the minima of empirical potentials are likely to generalize to those found by more-accurate computations, suggesting that the learned optimizers trained on these approximations will generalize as well. In this paper, we study a few functional forms of empirical potentials:

A.1. Lennard-Jones Clusters

First, the Lennard-Jones Clusters are the archetype of a simple-to-compute potential and are often used to model spherically-symmetric particles or atoms in free-space (a perfect vacuum with no other particles in the entire system). For example, this empirical potential can be used to model noble gasses or methane (Jones & Chapman, 1924). The energy landscape is defined only by pairwise distances between particles, denoted as $d_{ij}$ for atoms $i, j$.

$$\sum_i \sum_{j>i} \epsilon \left[ \left( \frac{d_0}{d_{ij}} \right)^{12} - 2 \left( \frac{d_0}{d_{ij}} \right)^{6} \right]$$

where $\epsilon$ describes the minimum two-particle energy and $d_0$ describes the distance where this occurs. Following prior work, in our paper we set $\epsilon$ and $d_0$ to 1, as the resulting minima structures can be scaled as necessary for systems where these settings do not hold.

While this model of atoms may seem simple, the corresponding optimization problem is anything but. For the corresponding task with 75 atoms, common gradient-based techniques such as Adam or FIRE cannot obtain the global minimum value and even the best trial out of 150 random initializations has an error greater than 7eV. Please see the main body of the paper for additional discussion of the difficulties of optimizing Lennard-Jones Clusters.

A.2. Gupta Clusters

The Gupta empirical potential adds an additional layer of complexity when modelling the energy of atomic structures. Designed to model lattice relaxations at a metal surface, the Gupta model provides an improved approximation by including a second-moment estimate of the tight-binding Hamiltonian (Gupta, 1981). The Gupta potential has been widely used in studying and predicting the stable structures of noble metals and bimetallic clusters.

The functional form of the Gupta potential is as follows:

$$\sum_i \sum_{j>i} A \exp \left[ p \left( 1 - \frac{d_{ij}}{d_0} \right) \right] - \xi \sum_i \sum_{j>i} \exp \left[ 2q \left( 1 - \frac{d_{ij}}{d_0} \right) \right]$$

The values $d_0, A, \xi, p, q$ are parameters of the Gupta potential that describe specific inter-particle interactions. In contrast to the Lennard-Jones system, these values cannot be factored out or set to defaults. Instead, the parameters are derived either experimentally or by fitting to Density-Functional Theory data from the bulk-faced cubic systems.

As the minima structures of the Gupta potential can vary based on the exact parameter values used, all of our experiments are based on the already-discovered configurations by Paz-Borbón et al. (2008). Table 4 provides the parameters used in our single-atom experiment modelling Au (gold) with 55 atoms and bimetallic experiments of Ag-Au, Ag-Pt, Pd-Au, and Pd-Pt clusters.

A.3. Stillinger-Weber (SW)

Stillinger-Weber potentials are designed to provide more accurate estimations of semiconductors and do so by including a three-body angular term (i.e. penalizing for deviations from an optimal angle within the crystal structure). In our paper, we use the Stillinger-Weber potential to model Silicon crystals. This system is distinct from the other two benchmarks, as the crystal model assumes the lattice structure is repeatedly infinitely in space (although a cutoff in interaction distance allows for models to use finite tilings).

For this system, the energy is defined by:
As the loss landscape of this system is ‘funneled’ and a low-energy paths exist between the local minima and the global minimum, we do not expect sporadic behavior of the learned optimizer to arise from the optimization problem itself. Instead, the erratic behavior of $\theta$ appears to be coming from instability in training, which is solved by our ESMC and GA training strategies.

### C. Baseline Optimizer Tuning

The baseline optimization techniques used in this paper (Adam, FIRE, Basin Hopping) have their own hyper-parameters which require tuning to obtain proper results. Doing so correctly is essential to ensure that our learned optimizer provide a meaningful improvement in minimum discovery that cannot be explained by improved tuning.

#### C.1. Tuning Adam and FIRE

In our paper, all baseline results arise from a two-stage process. First, we start with a grid search: utilizing 3 variants of the learning rate for each model, with values of \{0.01, 0.005, 0.001\}. Early exploration also modified the \(\beta\) parameters of Adam and the rate of increase/decrease of FIRE; however, both optimizers showed generic robustness to hyper-parameters other than learning rate changing.

While these baselines were somewhat competitive, we further tuned these optimizers by learning the values of the hyper-parameters for Adam and FIRE, following the Adam4p strategy by Metz et al. (2019b). This strategy, similar in spirit to the learned optimizers presented in the paper, uses meta-training to update the scalar parameters that define Adam and FIRE. 3 runs were conducted, initializing with each of the learning rates in the grid search. Meta-optimization was conducted for 100 outer-steps, applied with Adam with a learning rate of $10^{-2}$. The best hyper-parameters out of the grid search and the final learned parameters were then used to provide evaluation results.

#### C.2. Tuning Basin Hopping

Recall, Basin Hopping first uses standard optimization technique such as Adam or FIRE to optimize a network for a short number of steps. Parameters are perturbed from the minimum found, and optimization is performed once again to find a new minimum. If the new minimum is an improvement over the previous, then the new state is accepted.

### B. Comparison of Training Strategies

In order to compare the meta-training strategies of ES, ESMC (ours) and GA (ours), we focus on a simplified setup of the learned optimizers. Specifically, to remove the noise originating from the meta-training on diverse tasks, we focus on models trained only on the 13-atom Lennard-Jones clusters. As the model is trained only on 1 task, we only need to train for 900 meta-updates before both ESMC (ours) and GA (ours) appear to reach the global minimum on almost every single initialization. In contrast, traditional ES appears to be unstable, deviating greatly in the meta-training loss.
D. Training Details

For the sake of reproducible results, we provide additional details about learned optimizer training, including descriptions of how tasks are selected for the meta-batches and about what it means to produce a random initialization for particles.

D.1. Task Selection

During meta-training, estimates of the meta-gradients are produced by averaging over \( \sim 80 \) instantiations of atomic structure optimization problems, defined by both a random initialization (see Appendix D.2) and a corresponding empirical potential to minimize. Prior work on learned optimizers would refer to this set of problems as a task distribution and sample 80 tasks used to compute the a single meta-update. However, as the learned optimizers trained in this paper are still in the few-task regime, we instead default to sampling \([80/m]\) copies of all \( m \) tasks.

D.2. Harmonic Initialization

In the context of atomic structure optimization, purely random initializations (i.e. uniform over a pre-defined box size) are problematic as atoms that are too close to one another will have very large forces early in optimization. This can result in one atom being moved far away from the others. As most molecular dynamics simulations (including ours) use a cutoff distance for atomic interactions, future optimization steps are unlikely to update and recover this atom. As more particles often corresponds to lower energies, the resultant structure will be worse off than initializations where all particles are incorporated into the final structure.

A number of strategies have been proposed to stably initialize sets of atoms or particles. One strategy is to initialize uniformly over a large box, large-enough so that the particles are unlikely to be close but small enough that cutoffs are not a problem. In practice, we found this method difficult to tune when working with a variety of systems.

Instead, we utilize harmonic initialization. This strategy starts by randomly initializing coordinates in a small box, of size 3.0 Å for all Lennard-Jones models. Before atomic structural optimization begins, we first optimize an soft-sphere potential, which only penalizes per-particles distances when atoms are within pre-defined cutoff. Optimizing this intermediate function ensures that structural optimization does not have excessively large forces at the beginning of training. In free-space, this occurs by simply spreading particles apart.

A functional form of the soft-sphere potential is provided below:

\[
\sum_i \sum_{j>i} \begin{cases} 
0 & \text{if } d_{ij} > 0.1 \\
\epsilon (\frac{1-d_{ij})^\alpha}{\alpha} & \text{otherwise}
\end{cases}
\]

where in our formulation \( \epsilon = 1 \) and \( \alpha = 2 \). By default, we use 1000 optimization steps performed by gradient descent with a learning rate of \( 10^{-3} \). Note, the learned optimizers are not sensitive to changes in the harmonic step count or learning rate; both defaults were chosen to allow excess time for the soft-sphere potential to be minimized to \( \approx 0 \).

The same strategy is used for Gupta potentials, with the only difference being that the box is increased to edge length 4.0 Å to accommodate for the increased size of the atoms. For crystal optimization with the Stillinger-Weber potentials, our initialization respects the periodic boundary conditions, so particles are optimized within the pre-defined lattice.

D.3. Learned Optimizer Initialization

All weights of the neural network used to parameterize the learned optimizer are initialized via a LeCun Normal initialization (LeCun et al., 2012), following the default in the FLAX library (Heek et al., 2020), except for the output layer which is initialized to have output variance of 0. This default ensures that the learned optimizer at the beginning of training does not yield divergent trajectories.

For meta-training, the parameters that have the most significant impact on performance are the \( \alpha, \beta, \gamma \) used in the output parameterization. Best models used

\[ \alpha = 0.1 \quad \beta = 1 \quad \gamma = -3 \]

We hypothesize that this setup is stable when combined with the 0 output initialization, as steps start out very small and increase in size over the course of meta-training. To large of a step early in training may yield the undesirable scenario of particles becoming too close to one another.
Table 6. Transfer performance between potential types show that learned optimizers trained on Lennard-Jones can generalize to Stillinger-Weber, outperforming Adam and FIRE. However, the reverse is not true as the models trained on Silicon find very poor local minima, likely due to the removal of the periodic boundary condition.

| Training Set                  | Mean of 150 Initializations |
|-------------------------------|-----------------------------|
|                               | Lennard-Jones 13 atoms | Lennard-Jones 75 atoms | Stillinger-Weber Silicon 64 |
| **Baselines**                 |                            |                          |                           |
| Adam                          | -40.6                      | -380.5                   | -256.8                     |
| FIRE                          | -40.5                      | -380.2                   | -257.0                     |
| **Learned Optimizers**        |                            |                          |                             |
| Lennard-Jones {13, 19, 31, 38, 55, 75} | -44.3                      | -390.3                   | -261.7                     |
| Silicon 64                    | -35.6                      | -258.3                   | -261.8                     |
| **Global Minima**             | -44.3                      | -397.5                   | -277.2                     |

D.4. Compute Costs

Training costs vary significantly based on the system and distributed setup. For example, costs scale quadratically in the number of atoms for the Lennard-Jones and Gupta cluster and cubically for the Stillinger-Weber models due to the three-body terms. Rough estimates of the training time are 30 GPU hours for the Lennard-Jones models (single atom type) with the main bottleneck coming from the optimization of the 75 atom system. The Gupta models took about 10 GPU hours due to the smaller size.

E. Implementation

As mentioned in the core body of the text, the empirical potentials make use of JAX_MD (Schoenholz & Cubuk, 2019) and the optimization makes use of pure JAX (Bradbury et al., 2018). Code will be made available shortly.

F. Additional Results

In this section, we present additional results that were unable to fit in the main body of the paper. These results including minimal training task distributions and results beyond atomic structure optimization further support the use of learned optimizers in these global minimization problems.

F.1. Training with a Single Task

The results presented in Section 4 show significant benefits on global minima discovery when training only on a subset of 6 optimization tasks (defined by a different number of Lennard-Jones atoms). This diverse training set provides examples of both simple ‘funneled’ landscapes and glassy landscapes with large energy barriers between minima. However, it may be possible to perform well with significantly fewer training tasks.

In Figure 8, we show a comparison between training the learned optimizer on only Lennard-Jones with 13 atoms (LJ-13), Lennard-Jones with 75 atoms (LJ-75), and the diverse Lennard-Jones set from the main body of the paper {13, 19, 31, 38, 55, 75}. Interestingly, the models trained only on 13 or 75 atoms often perform better than Adam and FIRE and generalize significantly beyond the respective training distributions. The 13 atom model is perhaps the most impressive as it is the cheapest to train (due to the quadratic slowdown with number of atoms); but the model taking into account diverse examples show greater generalization to large atom counts.

F.2. Transfer Between Potentials

Early results in Table 6 also show that learned optimizers can transfer between empirical potentials; for example, the models trained on Lennard Jones clusters can transfer to the Stillinger-Weber potential despite differences in the periodic boundary and the addition of the angular term (but not vice versa). We believe that this form of generalization is most interesting and hope that future work explores this direction.
further; learned optimizers that train on empirical potential
and can be applied to DFT simulation appears a promising
avenue for significantly speeding up material design.