SINGULAR EQUIVARIANT ASYMPTOTICS AND WEYL’S LAW.
ON THE DISTRIBUTION OF EIGENVALUES OF AN
IN Variant ELLIPTIC OPERATOR

PABLO RAMACHER

Abstract. We study the spectrum of an invariant, elliptic, classical pseudodifferential operator
on a closed $G$-manifold $M$, where $G$ is a compact, connected Lie group acting effectively and
isometrically on $M$. Using resolution of singularities, we determine the asymptotic distribution
of eigenvalues along the isotypic components, and relate it with the reduction of the corresponding
Hamiltonian flow, proving that the reduced spectral counting function satisfies Weyl’s law,
together with an estimate for the remainder.

CONTENTS

1. Introduction 1
2. Fourier integral operators and equivariant asymptotics 5
3. Compact group actions and the momentum map 13
4. The generalized stationary phase theorem and resolution of singularities 16
5. The desingularization process 19
6. Phase analysis of the weak transforms. The first main theorem 32
7. Phase analysis of the weak transforms. The second main theorem 37
8. Asymptotics for the integrals $I_{\theta_1 \ldots \theta_N}(\mu)$ 42
9. Statement of the main result 46
10. On the spectrum of $\Gamma \backslash G$ 51
References 52

1. Introduction

The asymptotic distribution of eigenvalues of an elliptic operator has been object of mathematical
research for a long time. It was first studied by Weyl \cite{42} for certain second order differential
operators in Euclidean space using variational techniques, followed by work of Carleman \cite{11}, Min-
akashishundaram and Pleijel \cite{35}, Gårding \cite{20}, and Avacunović \cite{4}. Later, Hörmander \cite{27}
and Duistermaat-Guillemin \cite{16} extended these results to elliptic pseudodifferential operators on com-
 pact manifolds within the theory of Fourier integral operators. In this paper, we shall consider
this problem in the case that additional symmetries are present.

Let $M$ be a compact, connected, $n$-dimensional Riemannian manifold without boundary, $dM$
itits volume density, and

\[ P_0 : C^\infty(M) \rightarrow L^2(M) \]

an elliptic, classical pseudodifferential operator of order $m$ on $M$, regarded as an operator in the
Hilbert space $L^2(M)$ of square integrable functions on $M$ with respect to $dM$, its domain being

\[ \text{Date: August 11, 2011.} \]
the space \(C^\infty(M)\) of smooth functions on \(M\). Assume that \(P_0\) is positive and symmetric, which implies that \(P_0\) has a unique self-adjoint extension \(P\). Due to the compactness of \(M\), the spectrum of \(P\) is discrete. Consider now in addition a compact, connected Lie group \(G\), acting effectively and isometrically on \(M\), and assume that \(P\) commutes with the regular representation of \(G\) in \(L^2(M)\). In this situation, each eigenspace of \(P\) becomes a unitary \(G\)-module, and it is a natural question to ask about the distribution of the spectrum of \(P\) along the isotypic components of \(L^2(M)\) in the decomposition

\[ L^2(M) \simeq \bigoplus_{\chi \in \hat{G}} L^2(M)(\chi), \]

and the way it is related to the reduction of the corresponding Hamiltonian flow. It is described by the reduced spectral counting function \(N_\chi(\lambda) = d_\chi \sum_{t \leq \lambda} \text{mult}_\chi(t)\), where \(\text{mult}_\chi(t)\) denotes the multiplicity of the unitary irreducible representation \(\pi_\chi\) corresponding to the character \(\chi \in \hat{G}\) in the eigenspace \(E_t\) of \(P\) belonging to the eigenvalue \(t\). Let \(T^*M\) be the cotangent bundle of \(M\), \(p(x, \xi)\) the principal symbol of \(P_0\), and \(S^*M = \{(x, \xi) \in T^*M : p(x, \xi) = 1\}\). In his classical paper [27], Hörmander showed that the spectral counting function \(N(\lambda) = \sum_{t \leq \lambda} \dim E_t\) satisfies Weyl’s law

\[ N(\lambda) = \frac{\text{vol} S^*M}{n(2\pi)^n} \lambda^{n/m} + O(\lambda^{(n-1)/m}), \quad \lambda \to +\infty, \]

and it has been a long-standing open question whether a similar description for \(N_\chi(\lambda)\) can be achieved. While in the general case of effective group actions the leading term was obtained via heat kernel methods by Donnelly [14] and Brüning–Heintze [10], estimates for the remainder are not accessible via this approach. On the other hand, the derivation of remainder estimates within the framework of Fourier integral operators meets with serious difficulties when singular orbits are present, and until recently could only be obtained for finite group actions, or actions with orbits of the same dimension as in the work of Donnelly [14], Brüning–Heintze [10], Brüning [9], Helffer–Robert [23, 24], Guillemin–Uribe [22], and El-Houakmi–Helffer [19]. It was only in Ramacher [37] and Cassanas–Ramacher [12] that first partial results towards more general group actions were obtained within the setting of approximate spectral projections using resolution of singularities. The goal of this paper is to generalize this approach, and give an asymptotic description of \(N_\chi(\lambda)\) analogous to (1) for general effective group actions within the theory of Duistermaat, Guillemin, and Hörmander.

In order to explain the difficulties in a more detailed way, denote by \(Q = (P)^{1/m}\) the \(m\)-th root of \(P\) given by the spectral theorem, which is a classical pseudodifferential operator of order 1 with principal symbol \(q(x, \xi) = p(x, \xi)^{1/m}\). If \(0 < \lambda_1 \leq \lambda_2 \leq \ldots\) are the eigenvalues of \(P\) repeated according to their multiplicity, the eigenvalues of \(Q\) are \(\mu_j = (\lambda_j)^{1/m}\). Let \(\{dE^Q_\mu\}\) be the spectral resolution of \(Q\). The starting point of the method of Fourier integral operators is the Fourier transform of the spectral measure

\[ U(t) = \int e^{-it\mu} dE^Q_\mu = e^{-itQ}, \quad t \in \mathbb{R}, \]

which constitutes a one-parameter group of unitary operators in \(L^2(M)\). Although \(U(t)\) itself is not trace-class, it has a distribution trace given by the tempered distribution

\[ \text{tr} U(\cdot) : \mathcal{S}(\mathbb{R}) \ni \varphi \mapsto \int \sum_{j=1}^{\infty} e^{-it\mu_j} \varphi(t) dt = \sum_{j=1}^{\infty} \check{\varphi}(\mu_j) < \infty, \]
which is the Fourier transform of the spectral distribution

\[ \sigma(\mu) = \sum_{j=1}^{\infty} \delta(\mu - \mu_j). \]

An asymptotic description of the spectrum of \( P \) is then attained by studying the singularities of the distribution kernel of \( U(t) \) and of \( \text{tr} U(\cdot) \) for small \( |t| \). To be more precise, let \( \Omega_{1/2} \) denote the bundle of half-densities over \( M \), and \( U_{1/2} \) the operator which assigns to \( u_0 \in C^\infty(M, \Omega_{1/2}) \) the solution \( u \in C^\infty(\mathbb{R} \times M, \Omega_{1/2}) \) of the Cauchy problem

\[ (i^{-1} \partial_t + Q_{1/2}) u = 0, \quad u(0, x) = u_0(x), \]

where \( Q_{1/2} u = dM^{1/2} Q(u dM^{-1/2}) \). Then \( U_{1/2} : C^\infty(M, \Omega_{1/2}) \to C^\infty(\mathbb{R} \times M, \Omega_{1/2}) \) can be characterized globally as a Fourier integral operator with kernel \( \sigma(U) \) for suitable \( Q \), \( U \in \mathcal{S}(\mathbb{R}) \times \mathbb{R} \times \mathbb{R} \times M \). If the underlying group action is not free, \( \Omega \) is no longer a submersion, so that \( \Omega \) is not a smooth manifold. Therefore \( \sigma(U) \) fails to be a Fourier integral operator in general, so that, a priori, it is not clear how to describe its singularities by the method of Duistermaat, Guillemin,
and Hörmander. Ultimately, the difficulties arise from the necessity to understand the asymptotic behavior of oscillatory integrals of the form
\[ I(\mu) = \int_{T^*Y} \int_G e^{\mu \Phi(x,\xi,g)} a(gx, x, \xi, g) \, dg \, d(T^*Y)(x,\xi), \quad \mu \to +\infty, \]
via the stationary phase theorem, where \((\kappa, Y)\) are local coordinates on \(M\), \(d(T^*Y)(x,\xi)\) is the canonical volume density on \(T^*Y\), and \(dg\) the volume density on \(G\) with respect to some left invariant Riemannian metric on \(G\), while \(a \in C_c^\infty(Y \times T^*Y \times G)\) is an amplitude which might also depend on \(\mu\), and \(\Phi(x,\xi,g) = (\kappa(x) - \kappa(gx), \xi)\). For this, it would be necessary that the critical set of the phase function \(\Phi(x,\xi,g)\)

\[ \text{Crit}(\Phi) = \{(x,\xi,g) \in (\Omega \cap T^*Y) \times G : g \cdot (x,\xi) = (x,\xi)\} \]

were a smooth manifold, which, nevertheless, is only true for free group actions. In the case of general effective actions the stationary phase theorem can therefore not immediately be applied to the study of integrals of the type \(I(\mu)\), compare [9].

In this paper, we shall show how to overcome this obstacle by partially resolving the singularities of \(\mathcal{C} = \{(x,\xi,g) \in \Omega \times G : g \cdot (x,\xi) = (x,\xi)\}\), and applying the stationary phase principle in a suitable resolution space. This will be achieved by constructing a resolution of the set

\[ \mathcal{N} = \{(x,\xi,g) \in \mathcal{M} : gx = x\}, \quad \mathcal{M} = M \times G, \]

which is equivalent to a monomialization of its ideal sheaf \(I_\mathcal{N} \subset \mathcal{E}_\mathcal{M}\), where \(\mathcal{E}_\mathcal{M}\) denotes the structure sheaf of \(\mathcal{M}\). To be more precise, put \(X = T^*M \times G\), and let \(I_\mathcal{C} \subset \mathcal{E}_X\) be the ideal sheaf of \(\mathcal{C}\). Consider further the local ideal \(I_\Phi = (\Phi)\) generated by the phase function \(\Phi\), together with its vanishing set \(V_\Phi\). The derivative of \(I_\Phi\) is given by \(D(I_\Phi) = I_{\mathcal{C}|\mathcal{T}^*Y \times G}\), and \(\text{Crit}(\Phi) \subset V_\Phi\). The main idea is to construct a resolution of \(\mathcal{N}\), yielding a partial resolution \(\mathcal{Z} : \mathcal{X} \to \mathcal{C}\) of \(\mathcal{C}\), and a partial monomialization of \(I_\Phi\) according to

\[ \mathcal{Z}^*(I_\Phi) \cdot \mathcal{E}_{\mathcal{X},\mathcal{X}} = \prod_j \sigma_j^{l_j} \cdot \mathcal{Z}_-^{-1}(I_\Phi) \cdot \mathcal{E}_{\tilde{x},\mathcal{X}}, \quad \tilde{x} \in \mathcal{X}, \]
in such a way that \(D(\mathcal{Z}_-^{-1}(I_\Phi))\) is a resolved ideal sheaf. Here \(\mathcal{Z}^*(I_\Phi)\) denotes the inverse image ideal sheaf, \(\mathcal{Z}_-^{-1}(I_\Phi)\) the weak transform of \(I_\Phi\), while the \(\sigma_j\) are local coordinate functions, and \(l_j\) are natural numbers. As a consequence, the phase function factorizes locally according to \(\Phi \circ \mathcal{Z} \equiv \prod \sigma_j^{l_j} \cdot \mathcal{Z}^{(\cdot)}(\Phi)\), and we show that the weak transforms \(\mathcal{Z}^{(\cdot)}(\Phi)\) have clean critical sets in the sense of Bott [7]. An asymptotic description of the integrals \(I(\mu)\) can then be obtained by pulling them back to the resolution space \(\mathcal{X}\), and applying the stationary phase theorem to the weak transforms \(\mathcal{Z}^{(\cdot)}(\Phi)\) with the variables \(\sigma_j\) as parameters. The desingularization of \(\mathcal{N}\) will rely on the stratification of \(\mathcal{M}\) into orbit types, and consist of a series of monoidal transformations over \(\mathcal{M}\) where the centers are successively chosen as isotropy bundles over unions of maximally singular orbits.

The main result of the present paper is formulated in Theorem [12]. It states that the reduced spectral counting function satisfies Weyl’s law

\[ N_\chi(\lambda) = \frac{d_\chi[\pi_\chi:H]}{(n-\kappa)(2\pi)^{n-\kappa} \text{vol} ([\Omega \cap S^*M]/G)] \lambda^{n-\kappa} + O(\lambda^{(n-\kappa-1)/m} (\log \lambda)^\Lambda), \quad \lambda \to +\infty, \]
provided that \(n-\kappa \geq 1\), where \(\kappa\) is the dimension of a \(G\)-orbit of principal type, \(d_\chi\) the dimension of the irreducible representation \(\pi_\chi\), \([\pi_\chi:H] : 1\) the multiplicity of the trivial representation in the restriction of \(\pi_\chi\) to a principal isotropy group \(H\), and \(\Lambda\) a natural number which is bounded by the number of orbit types of the \(G\)-action on \(M\). The paper itself is structured as follows. Section 2 describes the theory of Duistermaat, Guillemin and Hörmander of spectral asymptotics in the
SINGULAR EQUIVARIANT ASYMPTOTICS AND WEYL’S LAW

5
equivariant setting, and explains how the problem of determining \( N_\chi (\lambda) \) reduces to the study of integrals of the type \( I(\mu) \) as \( \mu \to +\infty \). Section 3 contains some general remarks on compact group actions and the momentum map, followed by the computation of the critical set of the phase function \( \Phi \). Singular asymptotics are discussed in Section 4 after a brief account on the stationary phase principle and resolution of singularities. In Section 5 the desingularization process is carried out, giving way in Sections 6 and 7 to the phase analysis of the weak transforms \( \tilde{\Phi}^{wk} \). Asymptotics for integrals of the type \( I(\mu) \) are then obtained in Section 8 while the proof of the main result is given in Section 9.

Singular equivariant asymptotics with reminder estimates were previously obtained by Brüning-Heintze [10] and Duistermaat-Kolk-Varadarajan [18] for the spectrum of a discrete, uniform subgroup \( \Gamma \) of a connected, semisimple Lie group \( G \) with maximal compact subgroup \( K \). In the first case, a reminder estimate for the Gelfand-Gangolli-Wallach formula is given, which describes the distribution of eigenvalues of the Casimir operator along the isotypic components of \( L^2(\Gamma \setminus G) \). For torsion-free \( \Gamma \), this corresponds to the distribution of eigenvalues of the Bochner-Laplace operator on the spaces \( L^2(\Gamma \setminus G/ K, E^\chi) \), where \( E^\chi \) denotes the vector bundle on \( \Gamma \setminus G/ K \) induced by an arbitrary \( \chi \in \hat{K} \). In the second case, and under the assumption that \( \Gamma \) has no torsion, asymptotics for the spectral counting function of the Laplace-Beltrami operator \( \Delta \) on \( L^2(\Gamma \setminus G/ K) \) are derived. This amounts to an asymptotic description of \( N_\chi (\lambda) \) for \( \Delta \) on \( L^2(\Gamma \setminus G) \) in case that \( \chi \) corresponds to the trivial representation, and Theorem 12 generalizes this result to arbitrary \( \chi \in \hat{K} \), and subgroups \( \Gamma \) with torsion, as well as arbitrary invariant, elliptic, classical pseudodifferential operators. This is explained in Section 10.

Acknowledgments. The author wishes to express his gratitude to Mikhail Shubin for introducing him to this subject, and to his former collaborator Roch Cassanas. He also would like to thank Richard Melrose, Werner M"uller and Michèle Vergne for valuable conversations. This research was completed while the author was a member of the Mathematical Institute of G"ottingen University, and financed by the grant RA 1370/2-1 of the German Research Foundation (DFG).

2. FOURIER INTEGRAL OPERATORS AND EQUIVARIANT ASYMPTOTICS

Generalities. Let \( M \) be a compact, connected, \( n \)-dimensional Riemannian manifold, and \( G \) a compact, connected Lie group of dimension \( d \), acting effectively and isometrically on \( M \). Denote the canonical volume density on \( M \) by \( dM \) [40], page 112, and choose a left invariant Riemannian metric on \( G \) with volume density \( dg \). Let \( P_0 \) be an elliptic, classical pseudodifferential operator of order \( m \) on \( M \), regarded as an operator in \( L^2(M) \) with domain \( C^\infty(M) \), and assume that \( P_0 \) is positive and symmetric. Then \( P_0 \) has a unique self-adjoint extension \( P \) with the \( m \)-th Sobolev space \( H^m(M) \) as domain. Moreover, the spectrum of \( P \) is discrete. Assume now that \( P \) commutes with the regular representation of \( G \) in \( L^2(M) \) given by

\[
T(g)\varphi(x) = \varphi(g^{-1}x), \quad g \in G.
\]

Then every eigenspace of \( P \) becomes a unitary \( G \)-module, and it is natural to ask about the distribution of the spectrum of \( P \) along the isotypic components of \( L^2(M) \), which is described by the reduced spectral counting function \( N_\chi (\lambda) \) introduced in the previous section. We shall study this problem within the theory of Fourier integral operators developed by Hörmander, Duistermaat and Guillemin [27][16], and consider for this the \( m \)-th root \( Q = (P)^{1/m} \) of \( P \) given by the spectral theorem. By Seeley, \( Q \) is a classical pseudodifferential operator of order 1 with principal symbol \( q(x, \xi) = p(x, \xi)^{1/m} \) and domain \( H^1(M) \). If \( 0 < \lambda_1 \leq \lambda_2 \leq \ldots \) are the eigenvalues of \( P \) repeated according to their multiplicity, the eigenvalues of \( Q \) are \( \mu_j = (\lambda_j)^{1/m} \). Denote by \( \{dE^\chi_Q\} \) the spectral resolution of \( Q \). The starting point of the method developed by Hörmander, and which
goes back to work of Avacumović and Lewitan, is the Fourier transform of the spectral measure

\[ U(t) = \int e^{-it\mu}dE^Q_{\mu} = e^{-itQ}, \quad t \in \mathbb{R}, \]

which constitutes a one-parameter group of unitary operators in \( L^2(M) \). Now, if \( \{e_j\} \) denotes an orthonormal basis of eigenfunctions in \( L^2(M) \) of \( Q \) corresponding to the eigenvalues \( \{\mu_j\} \), then

\[ (2) \quad U(t)u = \sum_{j=1}^{\infty} e^{-it\mu_j} (u, e_j)_{L^2} e_j, \]

where \( u \in C^\infty(M) \), and \( u \in H^s(M), s \in \mathbb{Z} \), the sum converging in the \( C^\infty_- \), and \( H^s \)-topology, respectively, see \[33\], page 151. Thus, the distribution kernel of the operator \( U(t) : C^\infty(M) \to C^\infty(M) \subset D'(M) \) can be written as

\[ U(t,x,y) = \sum_{j=1}^{\infty} e^{-it\mu_j} e_j(x) e_j(y) \in D'(M \times M). \]

Although \( U(t) \) itself is not trace-class, it has a distribution trace given by the tempered distribution

\[ \text{tr} \ U(\cdot) : S(\mathbb{R}) \ni \varphi \mapsto \int \sum_{j=1}^{\infty} e^{-it\mu_j} \varphi(t)dt = \sum_{j=1}^{\infty} \hat{\varphi}(\mu_j) < \infty. \]

Indeed, for \( N_0 \in \mathbb{N} \), \( P^{-N_0} \) is a classical pseudodifferential operator of order \( -N_0m \). If \( N_0m > n \), its kernel is continuous, and \( P^{-N_0} \) is Hilbert-Schmidt, so that \( \sum_{j=1}^{\infty} \lambda_j^{-2N_0} < \infty \). Moreover, for \( \varphi \in S(\mathbb{R}) \) the infinite sum \( \sum_{j=1}^{\infty} \hat{\varphi}(\mu_j) e_j(x) \overline{e_j(y)} \) converges in \( C^\infty(M \times M) \), see \[21\], page 133. Because the Fourier transform is an isomorphism in \( S(\mathbb{R}) \), we conclude that \( \text{tr} \ U(t) = \sum_{j=1}^{\infty} e^{-it\mu_j} \) is the Fourier transform of the spectral distribution

\[ \sigma(\mu) = \sum_{j=1}^{\infty} \delta(\mu - \mu_j), \]

proving at the same time that \( \sigma \) is tempered. An asymptotic description of the spectrum of \( P \) is then attained by studying the singularities of \( U(t,x,y) \) and \( \text{tr} \ U(\cdot) \) for small \( |t| \). For this, Hörmander locally approximated the operator \( U(t) \) by Fourier integral operators, which solve the Cauchy problem approximately. More precisely, let \( U_{1/2} \) be the operator which assigns to \( u_0 \in C^\infty(M, \Omega_{1/2}) \) the solution \( u \in C^\infty(\mathbb{R} \times M, \Omega_{1/2}) \) of the hyperbolic Cauchy problem

\[ (i^{-1} \partial_t + Q_{1/2})u = 0, \quad u(0,x) = u_0(x), \]

where \( \Omega_{1/2} \) denotes the bundle of half-densities over \( M \), and \( Q_{1/2}u = dM_{1/2}^1 Q(u, dM_{-1/2}) \). It can then be shown \[16\], Theorem 1.1, that \( U_{1/2} : C^\infty(M, \Omega_{1/2}) \to C^\infty(\mathbb{R} \times M, \Omega_{1/2}) \) can be characterized globally as a Fourier integral operator with kernel \( \mathcal{U} \in I^{-1/4}(\mathbb{R} \times M, M, C^0) \) and canonical relation

\[ (3) \quad C = \{(t,\tau), (x, \xi), (y, \eta)) : (x, \xi), (y, \eta) \in T^*M \setminus 0, (t, \tau) \in T^*\mathbb{R} \setminus 0, \]

\[ \tau + q(x, \xi) = 0, (x, \xi) = \Phi^t(y, \eta) \}, \]

where \( \Phi^t \) is the flow in \( T^*M \setminus 0 \) of the Hamiltonian vector field associated to \( q \). This implies that the Fourier transform \( U_{1/2}(t) : C^\infty(M, \Omega_{1/2}) \to C^\infty(M, \Omega_{1/2}) \) of the spectral measure of \( Q_{1/2} \) is a Fourier integral operator of order 0 defined by the canonical transformation \( \Phi^t \), and that \( \hat{\sigma} \) can be characterized as a Fourier integral operator too, see \[16\], pp. 66. Moreover,

\[ \text{sing supp} \ U_{1/2} = \{(t, x, y) \in \mathbb{R} \times M \times M : (x, \xi) = \Phi^t(y, \eta) \text{ for suitable } \xi \in T^*_xM \setminus 0, \eta \in T^*_yM \setminus 0 \} \],
and similarly, $WF(\hat{\sigma}) \subset \{ (t, \tau) : \tau < 0 \text{ and } (x, \xi) = \Phi^t(x, \xi) \text{ for some } (x, \xi) \}$, so that $\hat{\sigma}$ is smooth on the complement of the set of periodic orbits. The study of the singularity of $\hat{\sigma} = \text{tr } U$ at $t = 0$ then leads to the main result of Hörmander

$$\hat{\sigma}(\hat{g}e^{i(t)\mu}) = \sum_{j=1}^{\infty} \hat{g}(\mu - \mu_j) \sim (2\pi)^{1-n} \sum_{k=0}^{\infty} c_k \mu^{n-1-k}, \quad \mu \to +\infty,$$

for suitable $g \in S(\mathbb{R})$, $\hat{g}(t) = g(-t)$, and with in principle known coefficients $c_k$, while for $\mu \to -\infty$ the expression is rapidly decreasing. From this, Weyl's classical law [11] follows by a Tauberian theorem.

Let us now come back to our initial question. To obtain a description of $N_\chi(\lambda)$, and to understand the way it is related to the reduction of the corresponding Hamiltonian flow [22], we would like to find an asymptotic expansion of $\sum_{j=1}^{\infty} m^{\mathcal{Q}}(\mu_j) \hat{g}(\mu - \mu_j)$ for suitable $g \in S(\mathbb{R})$, where $m^{\mathcal{Q}}(\mu_j) = d_\chi \text{mult}_\chi(\mu_j)/\text{dim } E^{\mathcal{Q}}_{\mu_j}$. This amounts to study the singularities of $\sum_{j=1}^{\infty} m^{\mathcal{Q}}(\mu_j) e^{-i\mu_j} \in S'(\mathbb{R})$. It corresponds to the distribution trace of $P_\chi \circ U(t)$, $P_\chi$ being the projector onto the $\chi$-isotypic component $L^2(M)(\chi)$, and is the Fourier transform of

$$\sigma_\chi(\mu) = \sum_{j=1}^{\infty} m^{\mathcal{Q}}(\mu_j) \hat{g}(\mu - \mu_j).$$

In what follows, denote by $\pi : \mathbb{R} \times G \times M \to \mathbb{R}$ the projection $(t, g, x) \mapsto t$, and by $\Gamma : \mathbb{R} \times G \times M \to \mathbb{R} \times M \times M$ the mapping $(t, g, x) \mapsto (t, x, gx)$. The global theory of Fourier integral operators [14], Lemma 7.1, implies that the transposed of the pullback, or pushforward $\pi_* : \mathcal{D}'(\mathbb{R} \times G \times M) \to \mathcal{D}'(\mathbb{R})$, can be characterized as a Fourier integral operator of class $\Gamma^{n/4-d/4}(\mathbb{R}, \mathbb{R} \times G \times M, C_1)$ with canonical relation

$$C_1 = \{(t, \tau) ; (t, \tau), (g, 0), (x, 0) \}.$$

It amounts to integration over $M \times G$. Similarly, the pullback $\Gamma^* : C^\infty(\mathbb{R} \times M \times M) \to C^\infty(\mathbb{R} \times G \times M)$ constitutes a Fourier integral operator of class $\Gamma^{n/4-d/4}(\mathbb{R} \times G \times M, \mathbb{R} \times M \times M, C_2)$ with canonical relation

$$C_2 = \{(t, \tau), (g, x^*\xi_2), (x, \xi_1 + g^*\xi_2) ; (t, \tau), (x, \xi_1), (gx, \xi_2) \},$$

where the map $x : G \to M$ is given by $g \mapsto gx$, and the map $g : M \to M$ by $x \mapsto gx$. A computation then shows that if we regard $\hat{\sigma}_\chi$ as a distribution density on $\mathbb{R}$ of order $1/2$, and $\pi_*$ and $\Gamma^*$ as maps between half densitites,

$$\hat{\sigma}_\chi = \sum_{j=1}^{\infty} m^{\mathcal{Q}}(\mu_j) e^{-i(\xi\mu_j)} = d_\chi \pi_* \hat{\chi} \Gamma^* U,$$

compare [16], page 66, and [14], Section 7. Now, although $\pi_*, \hat{\chi} \Gamma^*$, and $U_{1/2}$ are Fourier integral operators, their composition is not necessarily a Fourier integral operator. Indeed, the composition of the canonical relations of $\pi_*$ and $\Gamma^*$ reads

$$C_1 \circ C_2 = \{(t, \tau) ; (t, \tau), (x, \xi_1), (gx, \xi_2) \} : x^*\xi_2 = 0, \xi_1 + g^*\xi_2 = 0 \}.$$

But $x^*\xi_2 = 0$ means that $\xi_2 \in \text{Ann } T_x(G \cdot x)$. As will be explained in the next section, this is equivalent to $(x, \xi_2) \in \Omega = \mathbb{J}^{-1}(0)$, where $\mathbb{J} : T^*M \to \mathfrak{g}^*$ is the canonical symplectic momentum map, and we obtain

$$C_1 \circ C_2 \circ C = \{(t, \tau) : \text{ there exist } x, \eta, g \text{ such that } (x, \eta) \in \Omega, \quad (x, -g^*\eta) = \Phi^t(gx, \eta), \quad \tau + g(x, -g^*\eta) = 0 \}.$$

The singularities of $\hat{\sigma}_\chi$ are therefore determined by the restriction of $\Phi^t$ to $\Omega$. Since for general effective group actions, the zero level $\Omega$ is not smooth, neither $C_1 \circ C_2$ nor $C_1 \circ C_2 \circ C$ are smooth.
submanifolds in this case. Consequently, neither \( \pi_* \chi^* \Gamma^* \) nor \( \hat{\sigma}_\chi \) are Fourier integral operators in general. This faces us with serious difficulties when trying to study the singularities of \( \hat{\sigma}_\chi \) within the theory of Hörmander, Duistermaat, and Guillemin.

**A trace formula.** In what follows, we would like to understand the main singularity of \( \hat{\sigma}_\chi \) at \( t = 0 \) in greater detail. To this end, we shall first express

\[
\hat{\sigma}_\chi(\hat{g}e^{i(t)\mu}) = \sum_{j=1}^{\infty} m^Q(\mu_j) \hat{g}(\mu - \mu_j), \quad \hat{g} \in \mathcal{S}(\mathbb{R}),
\]

as the \( L^2 \)-trace of a certain operator. Observe that \( \sum_{j=1}^{\infty} \hat{g}(\mu_j)e_j(x)e_j(y) \in C^\infty(M \times M) \) is the Schwartz kernel of the bounded operator

\[
\int_{-\infty}^{+\infty} \hat{\theta}(t)U(t)dt : L^2(M) \to L^2(M),
\]

which is defined as a Bochner integral. It is of \( L^2 \)-trace class, since its kernel is square integrable over \( M \times M \). Therefore

\[
P_\chi \circ \int_{-\infty}^{+\infty} \hat{\theta}(t)U(t)dt = \int_{-\infty}^{+\infty} \hat{\theta}(t)P_\chi \circ U(t)dt
\]

must be of trace class, too, where

\[
P_\chi = d_\chi \int_G \chi(g)T(g) \, dg
\]

denotes the projector onto the isotypic component \( L^2(M)(\chi) \), and \( d_\chi \) the dimension of the irreducible representation corresponding to the character \( \chi \in \hat{G} \). We assert that kernel of the operator \([4]\) is given by \( \sum_{j=1}^{\infty} \hat{\theta}(\mu_j)P_\chi e_j(x)e_j(y) \in C^\infty(M \times M) \). Indeed, by choosing the eigenfunctions \( \{e_j\} \) according to the decomposition of the eigenspaces of \( Q \) into isotypic components, we can assume that \( P_\chi e_j = 0 \) if \( e_j \notin L^2(M)(\chi) \), and \( P_\chi e_j = e_j \) otherwise. By Sobolev’s inequality we have

\[
\|P_\chi e_j\|_{C^k} \leq \|e_j\|_{C^k} \leq c' \|e_j\|_{H^{k+n+1}} \leq c'' \|Q^{k+n+1}e_j\|_{L^2} \leq c''' \|e_j\|^{k+n+1},
\]

showing that \( \sum_j \hat{\theta}(\mu_j)P_\chi e_j(x)e_j(y) \) converges in \( C^\infty(M \times M) \), and with \([2]\) one computes

\[
\int_{-\infty}^{+\infty} \hat{\theta}(t)P_\chi \circ U(t)dt \, u(x) = \int_{-\infty}^{+\infty} \hat{\theta}(t) \sum_{j=1}^{\infty} e^{-it\mu_j} (u, e_j)_{L^2} P_\chi e_j(x)dt
\]

\[
= \int_M u(y) \sum_{j=1}^{\infty} \int_{-\infty}^{+\infty} \hat{\theta}(t) e^{-it\mu_j} P_\chi e_j(x)e_j(y) \, dt \, dM(y), \quad u \in L^2(M),
\]

everything being absolutely convergent. As a consequence,

\[
\text{tr} \int_{-\infty}^{+\infty} \hat{\theta}(t)P_\chi \circ U(t)dt = \sum_{j=1}^{\infty} \hat{\theta}(\mu_j)(P_\chi e_j, e_j)_{L^2} = \sum_{j=1}^{\infty} \hat{\theta}(\mu_j)m^Q(\mu_j) = \hat{\sigma}_\chi(\hat{g}),
\]

and we obtain the following L\(^2\)-trace formula, which was already derived in \([9]\).

**Lemma 1.** Let \( \hat{g} \in \mathcal{S}(\mathbb{R}) \). Then

\[
\hat{\sigma}_\chi(\hat{g}e^{i(t)\mu}) = \text{tr} \int_{-\infty}^{+\infty} \hat{\theta}(t)e^{it\mu}P_\chi \circ e^{-itQ}dt = \text{tr} d_\chi \int_{-\infty}^{+\infty} \int_G \hat{\theta}(t)e^{it\chi(g)}T(g) \, \hat{g} \circ e^{-itQ} \, dg \, dt.
\]
Let us now recall that \( U_{1/2} : C^\infty(M, \Omega_{1/2}) \to C^\infty(\mathbb{R} \times M, \Omega_{1/2}) \) can be characterized globally as a Fourier integral operator of class \( I^{-1/4}(\mathbb{R} \times M, M, C') \) with canonical relation given by \( [29] \). This means that for each coordinate patch \((\kappa, Y)\), and sufficiently small \( t \in (-\delta, \delta)\), the kernel of \( U_{1/2}(t) \) can be described locally as an oscillatory integral of the form

\[
\tilde{U}(t, \tilde{x}, \tilde{y}) = \int_{\mathbb{R}^n} e^{i\psi(t, \tilde{x}, \tilde{y})} a(t, \tilde{x}, \tilde{y}) d\tilde{y}
\]
on any compactum in \( Y \times Y \), where \( \tilde{x}, \tilde{y} \in \tilde{Y} = \kappa(Y) \subset \mathbb{R}^n \), and \( a \in S^0_{phg} \) is a classical symbol with \( a(0, \tilde{x}, \tilde{y}) = 1 \), while \( \psi(t, \tilde{x}, \tilde{y}) - \langle \tilde{y}, \eta \rangle \) is the defining phase function of \( C \) in the sense that

\[
C' = \{(t, \partial \psi / \partial t, \tilde{x}, \partial \psi / \partial \tilde{x}), (\partial \psi / \partial \eta, -\eta)\},
\]

see [29], page 254. Here we employed the notation \( d\eta = (2\pi)^{-n} d\eta \), \( d\eta \) being Lebesgue measure in \( \mathbb{R}^n \). Since \( \tau + q(x, \xi) = 0 \) on \( C \), and \( (\tilde{x}, \partial \psi / \partial \tilde{x}) = (\partial \psi / \partial \eta, \eta) \) for \( t = 0 \), we deduce \( \hat{d}_{\bar{x}, \eta} \psi(0, \tilde{x}, \eta) = d\tilde{x}, \tilde{y} \hat{\psi}(\tilde{x}, \eta) \), so that \( \psi \) is the solution of the Hamilton-Jacobi problem

\[
\frac{\partial \psi}{\partial t} + q(x, \frac{\partial \psi}{\partial x}) = 0, \quad \psi(0, \tilde{x}, \eta) = \langle \tilde{x}, \eta \rangle,
\]

\( \psi \) being homogeneous of degree 1. To construct an approximation of \( U(t) : L^2(M) \to L^2(M) \), let \( \{(\kappa_\gamma, Y_\gamma)\} \) be an atlas for \( M \), \( \{f_\gamma\} \) a corresponding partition of unity, and

\[
\hat{v}(\eta) = \int_{\mathbb{R}^n} e^{-i\langle \tilde{y}, \eta \rangle} v(\tilde{y}) \tilde{y} d\tilde{y}, \quad v \in C^\infty_c(\tilde{Y}_\gamma),
\]

the Fourier transform of \( v \). Write \( (\kappa_\gamma^{-1})^* dM = \beta_\gamma d\tilde{y} \), and denote by \( \tilde{U}_\gamma(t) \) the operator

\[
[\tilde{U}_\gamma(t)v](\tilde{x}) = \int_{\mathbb{R}^n} e^{i\psi_\gamma(t, \tilde{x}, \eta)} a_\gamma(t, \tilde{x}, \eta)v(\eta) d\eta,
\]

\( a_\gamma \) and \( \psi_\gamma \) being as described above, and set \( \tilde{U}_\gamma(t)u = [\tilde{U}_\gamma(t)(u \circ \kappa_\gamma^{-1})] \circ \kappa_\gamma, \ u \in C^\infty_c(Y_\gamma) \), so that we obtain the diagram

\[
\begin{array}{ccc}
C^\infty_c(Y_\gamma) & \xrightarrow{\tilde{U}_\gamma(t)} & C^\infty(Y_\gamma) \\
| & \kappa_\gamma \downarrow & \kappa_\gamma \\
C^\infty_c(\tilde{Y}_\gamma) & \xrightarrow{\tilde{U}_\gamma(t)} & C^\infty(\tilde{Y}_\gamma)
\end{array}
\]

Consider further test functions \( f_\gamma \in C^\infty_c(Y_\gamma) \) satisfying \( \tilde{f}_\gamma \equiv 1 \) on \( \text{supp} \ f_\gamma \), and define

\[
\tilde{U}(t) = \sum_\gamma F_\gamma \tilde{U}_\gamma(t) \tilde{F}_\gamma,
\]

where \( F_\gamma, \ \tilde{F}_\gamma \) denote the multiplication operators corresponding to \( f_\gamma \) and \( \tilde{f}_\gamma \), respectively. Then the result of Hörmander implies that

\[
(6) \quad R(t) = U(t) - \tilde{U}(t) \text{ is an operator with smooth kernel},
\]
In the last two expressions are oscillatory integrals with suitable regularizations. With (6) and

\[ \hat{A} \]

the stationary phase theorem, let us remark that since

\[ \hat{A} \]

the singularity of

\[ J \]

where

\[ q \]

where we wrote

\[ \gamma \]

is homogeneous in

\[ \eta \]

and

\[ \tilde{\chi} \]

Clearly,

\[ F \]

and

\[ \tilde{\psi} \]

one has the equality

\[ \tilde{\tilde{F}} \]

where

\[ d \]

where

\[ e \]

where

\[ n \]

one computes for

\[ P \]

Next, one computes for

\[ P \]

For

\[ \gamma \]

Corollary 1.

\[ \chi \]

\[ \eta \]

The singularity of

\[ \chi \]

of degree 1, Taylor expansion for small

\[ t \]

So far we have expressed \( \hat{\sigma}_\chi \) as an oscillatory integral. In order to study it by means of the stationary phase theorem, let us remark that since \( \psi \) is homogeneous in \( \eta \) of degree 1, Taylor expansion for small \( t \) gives

\[ \psi \]

where we wrote \( \gamma \) as a smooth function \( \tilde{\chi} \) which is homogeneous in \( \eta \) of degree 1 satisfying

\[ \gamma \]

Let us now define

\[ \mathcal{F} \]

Clearly, \( \mathcal{F} \) is rapidly decaying as a function in \( \tau \). More precisely, since \( a_\tau \in S^0_{\text{phbg}} \),

\[ |\mathcal{F}(\tau, \tilde{x}, \eta)| \leq C_N (1 + \tau^2)^{-N}, \quad \forall N \in \mathbb{N}, \tilde{x} \in Y_\gamma, \eta \in \mathbb{R}^n, \]

for some constant \( C_N > 0 \) which depends only on \( N \). Next note that \( \gamma \) is constant \( > 0 \) for all \( \tilde{x} \) and \( \omega \in S^{n-1} = \{ \eta \in \mathbb{R}^n : ||\eta|| = 1 \} \). There must therefore exist a constant \( C > 0 \) such that

\[ C|\eta| \geq \gamma(\tilde{x}, \eta) \geq \frac{1}{C}|\eta| \quad \forall \tilde{x} \in Y_\gamma, \eta \in \mathbb{R}^n, \]
which implies that for fixed $\mu$, $F_1(\mu - q_1(\tilde{x},\eta), \tilde{x}, \eta)$ is rapidly decaying in $\eta$. This yields a regularization of the oscillatory integral in the previous corollary, and we obtain

$$\hat{\sigma}_\chi(qe^{i(\mu)}u) = \frac{d\chi}{(2\pi)^n} \sum_{\gamma} \int_G \int_{T^*Y_\gamma} e^{i(c_{\gamma}(x) - c\gamma(\tilde{x},\eta),\eta)} \chi(g)f_\gamma(x)f_\gamma(\mu - q(x,\eta),\gamma(x,\eta)) \int_{T^*Y_\gamma} \left| \mu - q_1(\tilde{x},\eta)\right|^{N+M} \left| 1 - q_1(\tilde{x},\eta) \right|^M \nu_\gamma,$$

But even more is true. If we replace $\mu$ by $-\nu$, then $(\mu - q_\gamma(\tilde{x},\eta))^2 \geq 2\nu q_\gamma(\tilde{x},\eta) \geq 2\nu\eta/C$. From we therefore infer that $\hat{\sigma}_\chi(qe^{i(\mu)}u)$ is rapidly decreasing as $\mu \to -\infty$, reflecting the positivity of the spectrum. Assume now that $|1 - q_\gamma(\tilde{x},\eta/\mu)| \geq \text{const} > 0$. Then

$$|F(\mu - q_\gamma(\tilde{x},\eta), \tilde{x}, \eta)| \leq CN^{+M} \left| \frac{1}{\mu}\left| 1 - q_1(\tilde{x},\eta/\mu) \right|^{N+M} \right| \left| 1 - q_1(\tilde{x},\eta) \right|^M$$

for arbitrary $N, M \in \mathbb{N}$. Let therefore $0 \leq \alpha \in C^\infty_c((1/2,3/2))$ be such that $\alpha \equiv 1$ in a neighborhood of 1, so that

$$1 - \alpha(q_\gamma(\tilde{x},\eta/\mu)) \not\equiv 0 \quad \Rightarrow \quad |1 - q_\gamma(\tilde{x},\eta/\mu)| \geq \text{const} > 0.$$

Substituting $\eta = \mu\tilde{\eta}'$, we can rewrite $\hat{\sigma}_\chi(qe^{i(\mu)}u)$ as

$$\hat{\sigma}_\chi(qe^{i(\mu)}u) = \frac{\mu^n d\chi}{(2\pi)^n} \sum_{\gamma} \int_{-\delta}^{+\delta} \int_G \int_{T^*Y_\gamma} e^{i(t,\gamma(x),\eta)} \psi(t,\gamma(x),\eta) \chi(g)f_\gamma(x) \gamma(x,\eta))^2 \left| \mu - q_1(\tilde{x},\eta)\right|^{N+M} \left| 1 - q_1(\tilde{x},\eta) \right|^M$$

where all integrals are absolutely convergent. Now, since $\gamma(0,\tilde{x},\omega) = q_1(\tilde{x},\omega)$, there exists a constant $C > 0$ such that for sufficiently small $t \in (-\delta,\delta)$

$$C \geq \gamma(t,\tilde{x},\omega) \geq \frac{1}{C} \quad \forall \tilde{x} \in \gamma, \omega \in K,$$

$K$ being a compactum. By introducing the coordinates $\eta = R\omega$, $R > 0$, $\gamma(t,\kappa_\gamma(x),\omega) = 1$, one finally arrives at the following

**Proposition 1.** Let $\delta > 0$ be sufficiently small, and $\varphi \in C^\infty_c(-\delta,\delta)$. Then, as $\mu \to +\infty$,

$$\hat{\sigma}_\chi(qe^{i(\mu)}u) = \frac{\mu^n d\chi}{(2\pi)^n} \sum_{\gamma} \int_R \int_R e^{i(t,\gamma(x),\eta)} \chi(g)f_\gamma(x) \gamma(x,\eta)) \nu_\gamma(Rq(x,\omega)) d\gamma(S^*Y_\gamma)(x,\omega) \nu_\gamma,$$

up to terms of order $O(\mu^{-\infty})$, where $S^*Y_\gamma = \{(x,\omega) \in T^*Y_\gamma : \gamma(t,\kappa_\gamma(x),\omega) = 1\}$. Here $d\gamma(S^*Y_\gamma)(x,\omega)$ denotes the quotient of the volume density on $T^*Y_\gamma$ by Lebesgue measure in $\mathbb{R}$ with respect to $\gamma(t,\tilde{x},\omega)$. On the other hand, $\hat{\sigma}_\chi(qe^{i(\mu)}u)$ is rapidly decaying as $\mu \to -\infty$. 

Assume now that $\mu \geq 1$. To study the limit of $\hat{\sigma}_\chi(qe^{i(\mu)}u)$ as $\mu \to +\infty$, we shall apply the stationary phase principle to the $R\gamma$-integral first, and then to the integral over $G \times S^*Y_\gamma$. For the later phase analysis, it will be convenient to replace the integration over $G \times S^*Y_\gamma$ by an integration over $G \times T^*Y_\gamma$. Let us therefore note that since $\alpha \in C^\infty_c((1/2,3/2))$, \[ 1/2 \leq Rq(x,\omega) \leq 3/2 \quad \forall x \in Y_\gamma, \omega \in (S^*Y_\gamma)_x. \]
For sufficiently small $\delta$ we can therefore assume that the $R$-integration is over a compact intervall in $\mathbb{R}^+$. Let now $\sigma \in C_c^\infty(\mathbb{R})$ be a non-negative function with $\int \sigma(s) ds = 1$, and define $\Delta_{\varepsilon,r}(s) = \varepsilon^{-1}\sigma((s-r)/\varepsilon)$, $r \in \mathbb{R}$. Then

$$\Delta_{\varepsilon,r} \to \delta_r \quad \text{as} \quad \varepsilon \to 0$$

with respect to the weak topology in $E'(\mathbb{R})$. Using this approximation of the $\delta$-distribution and the theorem of Lebesgue on bounded convergence we obtain for $\hat{\sigma}$ the task of examining the asymptotic behavior of integrals of the form

$$\int_{\mathbb{R}^n} e^{i\mu x} \sigma(x) dx \to \int_{\mathbb{R}^n} e^{i\mu x} \sigma(x) dx \quad \text{as} \quad \varepsilon \to 0$$

$$\int_{\mathbb{R}^n} e^{i\mu x} \sigma(x) dx \to \int_{\mathbb{R}^n} e^{i\mu x} \sigma(x) dx \quad \text{as} \quad \varepsilon \to 0$$

To obtain a complete description, we are therefore left with

$$\kappa, Y \text{ via the generalized stationary phase theorem, where } (9) \Phi(\kappa, Y)$$

This will occupy us for the rest of this paper.

Theorem 1. Let $\varrho \in C_c^\infty(-\delta, \delta)$, and $\mu \geq 1$. For sufficiently small $\delta$ one has the asymptotic expansion

$$\hat{\sigma}_\varrho(\varrho e^{i(\mu)\varepsilon}) = \frac{\mu^{n-1} \varrho(0)}{(2\pi)^n} \sum_{\gamma} \int_G \int_{T^*Y_{\gamma}} e^{i\mu \varrho(x) - \kappa_\gamma(gx, \eta)} \bar{\chi}(g) \tilde{f}_\gamma(x) J_\gamma(g, x)$$

$$\Delta_{\varepsilon,1}(q(x, \eta)) d(T^*Y_{\gamma})(x, \eta) dg + O(\mu^{-2}),$$

where

$$O(\mu^{-2}) = C \mu^{-2} \sum_{|\beta| \leq 2} \sup_{R, t} \left| \partial^\beta_{R, t} \int_G \int_{T^*Y_{\gamma}} e^{i\mu \varrho(x) - \kappa_\gamma(gx, \eta)} \bar{\chi}(g) \tilde{f}_\gamma(x) \right|.$$

For $\mu \to -\infty$, the expression $\hat{\sigma}_\varrho(\varrho e^{i(\mu)\varepsilon})$ is rapidly decaying.

Proof. Since $(t, R) = (0, 1)$ is the only critical point of $t - R$, the assertion follows from the classical stationary phase theorem [21], Proposition 2.3.

We have thus partially unfolded the singularity of $\hat{\sigma}_\varrho$ at $t = 0$. Theorem 1 shows that its structure is more involved than in the non-equivariant setting, or in the case of finite group actions, compare [16], pp. 46, and [9], pp 92. To obtain a complete description, we are therefore left with the task of examining the asymptotic behavior of integrals of the form

$$I(\mu) = \int_{T^*Y} \int_G e^{i\Phi(x, \xi, g)} a(gx, x, \xi, g) dg d(T^*Y)(x, \xi), \quad \mu \to +\infty,$$

via the generalized stationary phase theorem, where $(\kappa, Y)$ are local coordinates on $M$, and $dg$ is the volume density of a left invariant metric on $G$, while $a \in C_c^\infty(Y \times T^*Y \times G)$ is an amplitude which might also depend on $\mu$, and

$$\Phi(x, \xi, g) = \langle \kappa(x) - \kappa(gx), \xi \rangle.$$
3. Compact group actions and the momentum map

**Compact group actions.** We commence this section by briefly recalling some basic facts about compact group actions that will be needed later. For a detailed exposition, we refer the reader to [8]. Let \( G \) be a compact Lie group acting locally smoothly on some \( n \)-dimensional \( C^\infty \)-manifold \( M \), and assume that the orbit space \( M/G \) is connected. Denote the stabilizer, or isotropy group, of a point \( x \in M \) by

\[
G_x = \{ g \in G : g \cdot x = x \}.
\]

The orbit of \( x \in M \) under the action of \( G \) will be denoted by \( G \cdot x \) or, alternatively, by \( O_x \), and is homeomorphic to \( G/G_x \). The equivalence class of an orbit \( O_x \) under equivariant homeomorphisms is called its orbit type, and the conjugacy class \( (G_x) \) of \( G_x \) in \( G \) its isotropy type. Now, if \( K_1 \) and \( K_2 \) are closed subgroups of \( G \), a partial ordering of orbit and isotropy types is given by

\[
\text{type } (G/K_1) \leq \text{type } (G/K_2) \iff (K_2) \subseteq (K_1) \iff K_2 \text{ is conjugated to a subgroup of } K_1.
\]

One of the main results in the theory of compact group actions is the following

**Theorem 2** (Principal orbit theorem). There exists a maximum orbit type \( G/H \) for \( G \) on \( M \). The union \( M(H) \) of orbits of type \( G/H \) is open and dense, and its image in \( M/G \) is connected.

**Proof.** See [8], Theorem IV.3.1. \( \square \)

Orbits of type \( G/H \) are called of principal type, and the corresponding isotropy groups are called principal. A principal isotropy group has the property that it is conjugated to a subgroup of each stabilizer of \( M \). Let \( K \subseteq G \) be a closed subgroup containing \( H \). An orbit of type \( G/K \) is called singular, if \( \dim K/H > 0 \), and exceptional, if \( K/H \) is finite and non-trivial, in which case \( \dim G/K = \dim G/H \), but \( \text{type } (G/K) \neq \text{type } (G/H) \). The following result says that there is a stratification of \( G \)-spaces into orbit types.

**Theorem 3.** Let \( G \) and \( M \) be as above, \( K \) a subgroup of \( G \), and denote the set of points on orbits of type \( G/K \) by \( M(K) \). Then \( M(K) \) is a topological manifold, which is locally closed. Furthermore, \( M(K) \) consists of orbits of type less than or equal to \( G/K \). The orbit map \( M(K) \to M(K)/G \) is a fiber bundle projection with fiber \( G/K \) and structure group \( N(K)/K \).

**Proof.** See [8], Theorem IV.3.3. \( \square \)

Let now \( M_\tau \) denote the union of non-principal orbits of dimension at most \( \tau \).

**Proposition 2.** If \( \kappa \) is the dimension of a principal orbit, then \( \dim M/G = n - \kappa \), and \( M_\tau \) is a closed set of dimension at most \( n - \kappa + \tau - 1 \).

**Proof.** See [8], Theorem IV.3.8. \( \square \)

Here the dimension of \( M_\tau \) is understood in the sense of general dimension theory. In what follows, we shall write \( \text{Sing } M = M - M(H) = M_\kappa \). Clearly,

\[
\text{Sing } M = M_0 \cup (M_1 - M_0) \cup (M_2 - M_1) \cup \cdots \cup (M_\kappa - M_{\kappa-1}),
\]

where \( M_i - M_{i-1} \) is precisely the union of non-principal orbits of dimension \( i \), and \( M_{-1} = \emptyset \), by definition. Note that

\[
M_i - M_{i-1} = \bigcup_j M(H_j^i), \quad H_j^i \subset G, \quad \dim G/H_j^i = i,
\]

is a disjoint union of topological manifolds of possibly different dimensions. Now, a crucial feature of smooth compact group actions is the existence of invariant tubular neighborhoods.
Theorem 4 (Invariant tubular neighborhood theorem). Assume that $G$ acts smoothly on $M$, and let $A$ be a closed $G$-invariant submanifold of $M$. Then $A$ has an invariant tubular neighborhood, that is, there exists a smooth $G$-vector bundle $\xi : E \to A$ on $A$ together with an equivariant diffeomorphism $\psi : E \to M$ onto an open neighborhood $W$ of $A$ such that the restriction of $\psi$ to the zero section of $\xi$ is the inclusion of $A$ in $M$.

Proof. See [8], Theorem VI.2.2. \hfill \square

Furthermore, by taking a $G$-invariant metric on $M$, $W$ can be identified via the exponential map with a neighborhood of the zero section in the normal bundle $\nu(A)$ of $A$. From now on, let $M$ be a closed, connected Riemannian manifold, and $G$ a connected compact Lie group acting on $M$ by isometries. Relying on the stratification of $M$ into orbit types, one can construct a $G$-invariant covering of $M$ as follows, compare [50], Theorem 4.20. Let $(H_1), \ldots, (H_L)$ denote the isotropy types of $M$, and arrange them in such a way that

$$(H_i) \geq (H_j) \Rightarrow i \leq j.$$ 

By Theorem 3 $M$ has a stratification into orbit types according to $M = M(H_1) \cup \cdots \cup M(H_L)$, and the principal orbit theorem implies that the set $M(H_1)$ is open and dense in $M$, while $M(H_1)$ is a closed, $G$-invariant submanifold. Denote by $\nu_1$ the normal $G$-vector bundle of $(H_1)$, and by $f_1 : \nu_1 \to M$ a $G$-invariant tubular neighbourhood of $M(H_1)$ in $M$. Take a $G$-invariant metric on $\nu_1$, and put

$$D_t(\nu_1) = \{v \in \nu_1 : \|v\| \leq t\}, \quad t > 0.$$ 

We then define the compact, $G$-invariant submanifold with boundary

$$M_2 = M - f_1(\bar{D}_{1/2}(\nu_1)),$$ 

on which the isotropy type $(H_1)$ no longer occurs, and endow it with a $G$-invariant Riemannian metric with product form in a $G$-invariant collar neighborhood of $\partial M_2$ in $M$. Consider now the union $M_2(H_2)$ of orbits in $M_2$ of type $G/H_2$, a compact $G$-invariant submanifold of $M_2$ with boundary, and let $f_2 : \nu_2 \to M_2$ be a $G$-invariant tubular neighbourhood of $M_2(H_2)$ in $M_2$, which exists due to the particular form of the metric on $M_2$. Taking a $G$-invariant metric on $\nu_2$, we define

$$M_3 = M_2 - f_2(\bar{D}_{1/2}(\nu_2)),$$

which constitutes a compact $G$-invariant submanifold with corners and isotropy types $(H_3), \ldots, (H_L)$. Continuing this way, one finally obtains the decomposition

$$M = f_1(D_{1/2}(\nu_1)) \cup \cdots \cup f_L(D_{1/2}(\nu_L)),$$

where we identified $f_L(D_{1/2}(\nu_L))$ with $M_L$, which leads to the covering

(11) $$M = f_1(\bar{D}_1(\nu_1)) \cup \cdots \cup f_L(\bar{D}_1(\nu_L)), \quad f_L(\bar{D}_1(\nu_L)) = M_L.$$ 

We introduce now the set

(12) $$\mathcal{N} = \{(x, g) \in \mathcal{M} : gx = x\}, \quad \mathcal{M} = M \times G,$$

which will play an important role later. If all isotropy groups of the $G$-action on $M$ have the same dimension, that is, if there are no singular orbits, $\mathcal{N}$ is a smooth manifold. Otherwise, $\mathcal{N}$ is singular, as can be seen from Theorem 1. Clearly, $\mathcal{N} = \bigcup_k \text{Iso } M(H_k)$, where $\text{Iso } M(H_k) \to M(H_k)$ denotes the isotropy bundle on $M(H_k)$, and by Proposition 2 we have

$$\dim \text{Iso } M(H_k) = \dim M(H_k) + \dim H_k \leq n - \kappa + \tau - 1 + \dim G - \tau = \dim M(H_L) - 1,$$

where $1 \leq k \leq L - 1$, and $\tau = \dim G/H_k$. The regular part $\text{Reg } \mathcal{N}$ is given by the union over all total spaces $\text{Iso } M(H_k)$ with non-singular isotropy type $(H_k)$, and is in general not dense in $\mathcal{N}$. \hfill \square
The momentum map. We shall now discuss the canonical symplectic momentum map of a closed, connected Riemannian manifold $M$ on which a connected, compact Lie group $G$ acts by isometries, and the way it is related to our problem. Consider the cotangent bundle $\pi : T^*M \to M$, as well as the tangent bundle $\tau : T^*M \to TM$, and define on $T^*M$ the Liouville form

$$\Theta(\xi) = \tau(\xi)[\pi^*(\xi)], \quad \xi \in T^*M.$$ 

We regard $T^*M$ as a symplectic manifold with symplectic form

$$\omega = d\Theta,$$

and define for any element $X$ in the Lie algebra $\mathfrak{g}$ of $G$ the function

$$J_X : T^*M \to \mathbb{R}, \quad \eta \mapsto \Theta(X)(\eta),$$

where $X$ denotes the fundamental vector field on $T^*M$, respectively $M$, generated by $X$. Note that $\Theta(X)(\eta) = \pi(X(\pi(\eta)))$. Indeed, put $\gamma(s) = e^{-sX} \cdot \eta$, $s \in (-\varepsilon, \varepsilon)$ for some $\varepsilon > 0$, so that $\gamma(0) = \eta$, $\gamma(0) = X. \eta$. Since $\pi(e^{-sX} \cdot \eta) = e^{-sX} \cdot \pi(\eta)$, one computes

$$\pi_X(\eta) = \frac{d}{ds} \pi_X(\gamma(s))_{|s=0} = \frac{d}{ds} e^{-sX} \cdot \pi(\eta)_{|s=0} = \tilde{\pi}_X(\eta).$$

Therefore

$$\Theta(\xi)(\eta) = \pi(X)[\pi^*(\eta)] = \eta(\tilde{X}_\eta),$$

as asserted. The function $J_X$ is linear in $X$, and due to the invariance of the Liouville form

$$L_X \Theta = dJ_X + i_X \omega = 0, \quad \forall X \in \mathfrak{g},$$

where $L$ denotes the Lie derivative. This means that $G$ acts on $T^*M$ in a Hamiltonian way. The corresponding symplectic momentum map is then given by

$$\mathcal{J} : T^*M \to \mathfrak{g}^*, \quad \mathcal{J}(\eta)(X) = J_X(\eta).$$

As explained in the previous section, we are interested in the asymptotic behavior of integrals of the form $(\Box)$, and would like to study them by means of the generalized stationary phase theorem, for which we have to compute the critical set of the phase function $\Phi(x, \xi, g)$. Let $(\kappa, Y)$ be local coordinates on $M$ as in $(\Box)$, and write $\kappa(x) = (\tilde{x}_1, \ldots, \tilde{x}_n)$, $\eta = \sum \xi_i(d\tilde{x}_i)x \in T^*_x X_Y$. One computes then for any $X \in \mathfrak{g}$

$$\frac{d}{dt} \Phi(x, \xi, e^{tX})_{|t=0} = \frac{d}{dt} \langle \kappa(\xi - tX), x \rangle_{|t=0} = \sum \xi_i \tilde{x}_i(\tilde{x}_i) = \sum \xi_i(\tilde{x}_i)(\tilde{x}_i) = \eta(\tilde{x}_i) = \Theta(\tilde{X}_\eta)(\eta) = \mathcal{J}(\eta)(X).$$

Therefore $\Phi$ represents the global analogue of the momentum map; furthermore, their critical sets are essentially the same. Indeed, one has

$$\partial \xi \Phi(\kappa^{-1}(\tilde{x}), \xi, g) = [1 - T(\kappa \circ g \circ \kappa^{-1})]_x \xi = (1 - g_x^*) \xi,$$

so that $\partial_x \Phi(x, \xi, g) = 0$ amounts precisely to the condition $g^* \xi = \xi$. Since $\partial_x \Phi(x, \xi, g) = 0$ if, and only if $g_x = x$, one obtains

$$\text{Crit}(\Phi) = \{(x, \xi, g) \in T^*Y \times G : (\Phi)_*(x, \xi, g) = 0\} = \{(x, \xi, g) \in (\Omega \cap T^*Y) \times G : g \cdot (x, \xi) = (x, \xi)\},$$

where $\Omega = \mathcal{J}^{-1}(0)$ is the zero level of the momentum map. Note that

$$\eta \in \Omega \cap T^*_x Y \iff \eta \in \text{Ann}(T_x(G \cdot x)),$$

where $\text{Ann}(V_x) \subset T^*_x M$ denotes the annihilator of a vector subspace $V_x \subset T_x M$. Now, the major difficulty in applying the generalized stationary phase theorem in our setting stems from the fact that, due to the orbit structure of the underlying group action, the zero level $\Omega$ of the momentum map, and, consequently, the considered critical set $\text{Crit}(\Phi)$, are in general singular varieties. In fact,
if the $G$-action on $T^*M$ is not free, the considered momentum map is no longer a submersion, so that $\Omega$ and the symplectic quotient $\Omega/G$ are no longer smooth. Nevertheless, it can be shown that these spaces have Whitney stratifications into smooth submanifolds, see [39], and [36], Theorems 8.3.1 and 8.3.2, which correspond to the stratifications of $T^*M$, and $M$ by orbit types [17]. In particular, if $(H_L)$ denotes the principal isotropy type of the $G$-action in $M$, $\Omega$ has a principal stratum given by

$$\text{Reg } \Omega = \{ \eta \in \Omega : G_\eta \sim H_L \},$$

where $G_\eta$ denotes the isotropy group of $\eta$. To see this, let $\eta \in \Omega \cap T^*_x M$, and $G_x \sim H_L$. In view of (13) one computes for $g \in G_x$, and $x = x_F + x_N \in T_x M = T_x(G \cdot x) \oplus N_x(G \cdot x)$

$$g \cdot \eta(x) = \eta((L_g^{-1})_{*x}(x_N)) = \eta(x),$$

since $G_x$ acts trivially on $N_x(G \cdot x)$, see [8], pages 308 and 181. But $G_\eta \subset G_{\pi(\eta)}$ for arbitrary $\eta \in T^*M$, so that we conclude

$$\eta \in \Omega \cap T^*_x M, \ G_x \sim H_L \Rightarrow G_\eta = G_x.$$

Since the stratum Reg $\Omega$ is open and dense in $\Omega$, equality (15) follows. Note that Reg $\Omega$ is a smooth submanifold in $T^*M$ of codimension equal to the dimension $\kappa$ of a principal $G$-orbit in $M$. It is therefore clear that the smooth part of Crit($\Psi$) corresponds to

$$\text{Reg } \text{Crit}(\Psi) = \{ (x, \xi, g) \in (\text{Reg } \Omega \cap T^*Y) \times G : g \in G_{(x, \xi)} \},$$

and constitutes a submanifold of codimension $2\kappa$.

4. The generalized stationary phase theorem and resolution of singularities

**The principle of the stationary phase.** Since the critical set of the phase function (13) is not necessarily smooth, the stationary phase method can not immediately be applied to derive asymptotics for the integral (9). We shall therefore first partially resolve the singularities of Crit($\Psi$), and then apply the stationary phase principle in a suitable resolution space. To explain our approach, let us begin by recalling

**Theorem 5** (Generalized stationary phase theorem for manifolds). Let $M$ be a $n$-dimensional Riemannian manifold with volume density $dM$, $\psi \in C^\infty(M)$ a real valued phase function, and set

$$\mathcal{I}(\mu) = \int_M e^{i\psi(m)}/\mu a(m) \, dM(m), \quad \mu > 0,$$

where $a(m) \in C^\infty_c(M)$. Let

$$\mathcal{C} = \{ m \in M : \psi_\ast : T_m M \rightarrow T_{\psi(m)} \mathbb{R} \text{ is zero} \}$$

be the critical set of the phase function $\psi$, and assume that it is clean in the sense of Bott [7], meaning that

(I) $\mathcal{C}$ is a smooth submanifold of $M$ of dimension $p$ in a neighborhood of the support of $a$;

(II) at each point $m \in \mathcal{C}$, the Hessian $\psi''(m)$ of $\psi$ is transversally non-degenerate, i.e. non-degenerate on $T_m M/T_m \mathcal{C} \simeq N_m \mathcal{C}$, where $N_m \mathcal{C}$ denotes the normal space to $\mathcal{C}$ at $m$.

Then, for all $N \in \mathbb{N}$, there exists a constant $C_{N, \psi} > 0$ such that

$$|\mathcal{I}(\mu) - e^{i\psi_0(2\pi \mu)}(2\pi \mu)^{-\frac{m}{2}} \sum_{j=0}^{N-1} \mu^j Q_j(\psi; a)| \leq C_{N, \psi} \mu^N \text{vol}(\text{supp } a \cap \mathcal{C}) \sup_{l \leq 2N} \| D^l a \|_\infty, $$

where

$$\sum_{j=0}^{N-1} \mu^j Q_j(\psi; a) = \sum_{j=0}^{N-1} \mu^j \int_M e^{i\psi(m)/\mu} \cdot \frac{\partial^{2j} a(m)}{\partial \nu(m)^{2j}}(m) \, dM(m),$$

with $\nu(m)$ the inward normal of $\partial M$ at $m$.
where $D^l$ is a differential operator on $M$ of order $l$, and $\psi_0$ is the constant value of $\psi$ on $C$. Furthermore, for each $j$ there exists a constant $C_{j,\psi} > 0$ such that

$$|Q_j(\psi; a)| \leq \bar{C}_{j,\psi} \text{vol}(\text{supp } a \cap C) \sup_{l \leq 2j} \|D^l a\|_{\infty, C},$$

and, in particular,

$$Q_0(\psi; a) = \int_C \frac{a(m)}{|\det \psi''(m)|_{N_m C}|^{1/2}} d\sigma_C(m) e^{i \frac{2}{\pi} \psi''},$$

where $d\sigma_C$ is the induced volume density on $C$, and $\mu''$ the constant value of the signature of the transversal Hessian $\psi''(m)_{N_m C}$ on $C$.

**Proof.** See for instance [28], Theorem 7.7.5, together with [13], Theorem 3.3, as well as [41], Theorem 2.12. □

**Remark 1.** An examination of the proof of the foregoing theorem shows that the constants $C_{N,\psi}$ are essentially bounded from above by

$$\sup_{m \in C \cap \text{supp } a} \left\| \left( \psi''(m)_{N_m C} \right)^{-1} \right\|.$$

Indeed, let $\alpha : (x, y) \to m \in \mathcal{O} \subset M$ be local normal coordinates such that $\alpha(x, y) \in C$ if, and only if, $y = 0$. The transversal Hessian $\text{Hess} \psi(m)_{N_m C}$ is given in these coordinates by the matrix

$$\left( \partial_{y_k} \partial_{y_l} (\psi \circ \alpha)(x, 0) \right)_{k,l}$$

where $m = \alpha(x, 0)$, compare [28]. If the transversal Hessian of $\psi$ is non-degenerate at the point $m = \alpha(x, 0)$, then $y = 0$ is a non-degenerate critical point of the function $y \mapsto (\psi \circ \alpha)(x, y)$, and therefore an isolated critical point by the lemma of Morse. As a consequence,

$$\left| \frac{\partial y}{\partial \alpha}(\psi \circ \alpha)(x, y) \right| \leq 2 \left\| \left( \partial_{y_k} \partial_{y_l} (\psi \circ \alpha)(x, 0) \right)_{k,l} \right\|$$

for $y$ close to zero. The assertion now follows by applying [28], Theorem 7.7.5, to the integral

$$\int_{\alpha^{-1}(\mathcal{O})} e^{i \mu(x, y)/\mu} (\alpha \circ \alpha)(x, y) dy dx$$

in the variable $y$ with $x$ as a parameter, since in our situation the constant $C$ occurring in [28], equation (7.7.12), is precisely bounded by [13], if we assume as we may that $\alpha$ is supported near $C$. A similar observation holds with respect to the constants $\bar{C}_{j,\psi}$.

Conditions (I) and (II) in Theorem 5 are essential. Actually, the existence of singularities might alter the asymptotics, as can be seen from the following

**Example 1.** Let $M = \mathbb{R}^2$, $\psi(x, y) = (xy)^2$, and consider the asymptotic behavior of the integral

$$\mathcal{I}(\mu) = \int \int e^{i \psi(x, y)/\mu} a(x, y) dx dy$$

as $\mu \to 0^+$, where $a(x, y) \in C_\infty^c(\mathbb{R}^2)$ is a compactly supported amplitude, and $dx dy$ denotes Lebesgue measure in $\mathbb{R}^2$. The critical set of $\psi$ is given by the singular variety $\text{Crit}(\psi) = \{xy = 0\}$, and a computation shows that

$$\mathcal{I}(\mu) = \frac{e^{i \pi/4}}{\sqrt{2}} a(0, 0)(2\pi \mu)^{1/2} \log(\mu^{-1}) + O(\mu^{1/2}).$$

In general, one faces serious difficulties in describing the asymptotic behavior of integrals of the form [18] if the critical set $C$ is not smooth and in what follows, we shall indicate how to circumvent this obstacle by using resolution of singularities.
Resolution of singularities. Let $M$ be a smooth variety over a field of characteristic zero, $\mathcal{O}_M$ the structure sheaf of rings of $M$, and $I \subset \mathcal{O}_M$ an ideal sheaf. The aim of resolution of singularities is to construct a birational morphism $\Pi : \tilde{M} \to M$ such that $\tilde{M}$ is smooth, and the inverse image ideal sheaf $\Pi^*(I) \subset \mathcal{O}_{\tilde{M}}$, which is the ideal sheaf generated by the pullbacks of local sections of $I$, is locally principal. This is called the principalization of $I$, and implies resolution of singularities. That is, for every quasi-projective variety $X$, there is a smooth variety $\tilde{X}$, and a birational and projective morphism $\pi : \tilde{X} \to X$. Vice versa, resolution of singularities implies principalization. If $\Pi^*(I)$ is monomial, that is, if for every $x \in \tilde{M}$ there are local coordinates $z_i$ and natural numbers $c_i$ such that

$$\Pi^*(I) \cdot \mathcal{O}_{x,\tilde{M}} = \prod_i z_i^{c_i} \cdot \mathcal{O}_{x,\tilde{M}},$$

one obtains strong resolution of singularities, which means that, in addition to the properties stated above, $\pi$ is an isomorphism over the smooth locus of $X$, and $\pi^{-1}(\text{Sing } X)$ a divisor with simple normal crossings. By the work of Hironaka [26], resolutions are known to exist, and we refer the reader to [31] for a detailed exposition. Let next $D(I)$ be the derivative of $I$, which is the ideal sheaf that is generated by all derivatives of elements of $I$. Let further $Z \subset M$ be a smooth subvariety, and $\pi : B_Z M \to M$ the corresponding monoidal transformation with center $Z$ and exceptional divisor $F \subset B_Z M$. Assume that $(I, m)$ is a marked ideal sheaf with $m \leq \text{ord}_Z I$. The total transform $\pi^*(I)$ vanishes along $F$ with multiplicity $\text{ord}_Z I$, and by removing the ideal sheaf $\mathcal{O}_{B_Z M}(m \cdot F)$ from $\pi^*(I)$ we obtain the birational, or weak transform $\pi_*^{-1}(I, m) = (\mathcal{O}_{B_Z M}(mF) \cdot \pi^*(I), m)$ of $(I, m)$. Take now local coordinates $(x_1, \ldots, x_n)$ on $M$ such that $Z = (x_1 = \cdots = x_r = 0)$. As a consequence,

$$y_1 = \frac{x_1}{x_r}, \ldots, y_{r-1} = \frac{x_{r-1}}{x_r}, y_r = x_r, \ldots, y_n = x_n$$

define local coordinates on $B_Z \tilde{M}$, and for $(f, m) \in (I, m)$ one puts

$$\pi_*^{-1}(f(x_1, \ldots, x_n), n) = (y_r^{-m}\cdot f(y_1y_r, \ldots, y_{r-1}y_r, y_r, \ldots, y_n), m).$$

By computing the first derivatives of $\pi_*^{-1}(f(x_1, \ldots, x_n), m)$, one then sees that for any composition $\Pi : \tilde{M} \to M$ of blowing-ups of order greater or equal than $m$,

$$(20) \quad \Pi_*^{-1}(D(I, m)) \subset D(\Pi_*^{-1}(I, m)),$$

see [31], Sections 3.5 and 3.7.

Consider now an oscillatory integral of the form (18), and its asymptotic behavior as $\mu \to +0$, in case that the critical set $C$ of the phase function $\psi$ is not clean. The essential idea behind our approach to singular asymptotics via resolution of singularities is to obtain a partial monomialization

$$\Pi^*(I_\psi) \cdot \mathcal{O}_{x,\tilde{M}} = z_1^{c_1} \cdots z_k^{c_k} \Pi_*^{-1}(I_\psi) \cdot \mathcal{O}_{x,\tilde{M}}$$

of the ideal sheaf $I_\psi = (\psi)$ generated by the phase function $\psi$ via a suitable resolution $\Pi : \tilde{M} \to M$ in such a way that the corresponding weak transforms $\tilde{\psi}^{wk} = \Pi_*^{-1}(\psi)$ have clean critical sets in the sense of Bott [7]. Here $z_1, \ldots, z_k$ are local variables near each $x \in \tilde{M}$ and $c_i$ are natural numbers. This enables one to apply the stationary phase theorem in the resolution space $\tilde{M}$ to the weak transforms $\tilde{\psi}^{wk}$ with the variables $z_1, \ldots, z_k$ as parameters. Note that by Hironaka’s theorem, $I_\psi$ can always be monomialized. But in general, this monomialization would not be explicit enough to allow an application of the stationary phase theorem.

In the situation of the previous sections, consider the set $N$ defined in (12). To derive asymptotics for the integral (9), we shall construct a strong resolution of $N$, from which we shall deduce a partial desingularization $Z : \tilde{X} \to X = T^*M \times G$ of the set

$$(21) \quad C = \{(x, \xi, g) \in \Omega \times G : g \cdot (x, \xi) = (x, \xi)\},$$
and a partial monomialization of the local ideal \( I_\Phi = (\Phi) \) generated by the phase function \( \Phi \),

\[
Z^*(I_\Phi) \cdot E_{\tilde{x}, \tilde{X}} = \prod_j \sigma_j^{l_j} \cdot Z_{\sigma_j^{-1}}(I_\Phi) \cdot E_{\tilde{x}, \tilde{X}},
\]

where \( \sigma_j \) are local coordinate functions near each \( \tilde{x} \in \tilde{X} \), and \( l_j \) natural numbers. As a consequence, the phase function factorizes locally according to \( \Phi \circ Z = \prod \sigma_j^{l_j} \cdot \tilde{\Phi}^{wk} \), and we show that the weak transforms \( \tilde{\Phi}^{wk} \) have clean critical sets. Asymptotics for the integrals \( I(\mu) \) are then obtained by pulling them back to the resolution space \( \tilde{X} \), and applying the stationary phase theorem to the \( \tilde{\Phi}^{wk} \) with the variables \( \sigma_j \) as parameters.

A general description of the asymptotic behavior of oscillatory integrals with singular critical sets was given in [3], and later also in [15, 32, 2], using Hironaka’s theorem on resolution of singularities. It implies that integrals of the form (18) always have local expansions of the form

\[
\sum_{\alpha} \sum_{k=0}^{n-1} c_{\alpha k}(a) \mu^\alpha (\log \mu^{-1})^k, \quad \mu \to +0,
\]

where the coefficient \( \alpha \) runs through a finite set of arithmetic progressions of rational numbers, and the \( c_{\alpha k} \) are distributions on \( M \) with support in \( \mathcal{C} \). The occurring coefficients \( \alpha \) and \( k \) are determined by the so-called numerical data of the resolution, and their computation is in general a difficult task, unless one constructs an explicit resolution. Resolution of singularities was first employed in [6, 3] to give a new proof of the Hörmander-Lojasiewicz theorem on the division of distributions and hence to the existence of temperate fundamental solutions for constant coefficient differential operators. Since many problems in analysis originate in the singularities of some critical variety, it seems likely that an application of resolution of singularities may be relevant in further areas of this field.

Partial desingularizations of the zero level set \( \Omega \) of the moment map and the symplectic quotient \( \Omega / G \) have been obtained e.g. in [33] for compact symplectic manifolds with a Hamiltonian compact Lie group action by performing blowing-ups along minimal symplectic suborbifolds containing the strata of maximal depth in \( \Omega \). Recently, resolutions of group actions were also considered in [1] to study the equivariant cohomology of compact \( G \)-manifolds.

5. The desingularization process

We shall now proceed to resolve the singularities of (12). For this, we will have to set up an iterative desingularization process along the strata of the underlying \( G \)-action, where each step in our iteration will consist of a decomposition, a monoidal transformation, and a reduction. The centers of the monoidal transformations are successively chosen as isotropy bundles over unions of maximally singular orbits. For simplicity, we shall assume that at each iteration step the union of maximally singular orbits is connected. Otherwise each of the connected components, which might even have different dimensions, has to be treated separately.

First decomposition. Let \( M \) be a closed, connected Riemannian manifold, and \( G \) a connected compact Lie group acting on \( M \) by isometries. As in the previous section, let \( (H_1), \ldots, (H_L) \) be the isotropy types of the \( G \)-action on \( M \), \( 1 \leq k \leq L - 1 \), and \( f_k : \nu_k \to M_k \) an invariant tubular neighborhood of \( M_k(H_k) \) in

\[
M_k = M - \bigcup_{i=1}^{k-1} f_i(D_{1/2}^{\mathcal{C}}(\nu_i)),
\]

a manifold with corners on which \( G \) acts with the isotropy types \( (H_k), (H_{k+1}), \ldots, (H_L) \). Here

\[
f_k(p^{(k)}, v^{(k)}) = (\exp_{p^{(k)}} \circ \gamma^{(k)})(v^{(k)}), \quad p^{(k)} \in M_k(H_k), \quad v^{(k)} \in (\nu_k)_{p^{(k)}},
\]
is an equivariant diffeomorphism, while

$$\gamma^{(k)}(v^{(k)}) = \frac{F_k(p^{(k)})}{(1 + ||v^{(k)}||)^{1/2}} v^{(k)},$$

where $F_k : M_k(H_k) \to \mathbb{R}$ is a smooth, $G$-invariant, positive function, see [8], pp. 306. Let $S_k = \{ v \in \nu_k : ||v|| = 1 \} \to M_k(H_k)$ be the sphere bundle over $M_k(H_k)$, and put $W_k = f_k(D_1(\nu_k))$, $W_L = M_L$, so that

$$M = W_1 \cup \cdots \cup W_L,$$

see [11]. Endow $G$ with the Riemannian structure

$$d(X_g, Y_g) = -\text{tr} \, (dL_{g^{-1}}(X_g)) \, (dL_{g^{-1}}(Y_g)), \quad X_g, Y_g \in T_g G,$$

where $L_g : h \to gh, h \in G$, and introduce for each $p^{(k)} \in M_k(H_k)$ the decomposition

$$T_e G \simeq \mathfrak{g} = \mathfrak{g}_{p^{(k)}} \oplus \mathfrak{g}_{p^{(k)}}^\perp,$$

where $\mathfrak{g}_{p^{(k)}} \simeq T_e G_{p^{(k)}}$ denotes the Lie algebra of the stabilizer $G_{p^{(k)}}$ of $p^{(k)}$, and $\mathfrak{g}_{p^{(k)}}^\perp$ its orthogonal complement with respect to the above Riemannian structure. Note that $T_e G_{p^{(k)}} \simeq dL_{h}(\mathfrak{g}_{p^{(k)}})$, and if $A \in \mathfrak{g}_{p^{(k)}}^\perp$, $d(dL_{h}(X), dL_{h}(A)) = -\text{tr} \, (X) \, (dA) = 0$ for all $X \in \mathfrak{g}_{p^{(k)}}$. Therefore, the mapping

$$\mathfrak{g}_{p^{(k)}}^\perp \ni A \mapsto dL_{h}(A) = \frac{d}{dt}(he^{tA})|_{t=0} \in N_{h} G_{p^{(k)}}$$

establishes an isomorphism $\mathfrak{g}_{p^{(k)}}^\perp \simeq N_{h} G_{p^{(k)}}$. In fact, $\text{Ad} \,(G_{p^{(k)}}) \, \mathfrak{g}_{p^{(k)}}^\perp \subset \mathfrak{g}_{p^{(k)}}^\perp$, so that $G / G_{p^{(k)}}$ constitutes a reductive homogeneous space, while the distribution $G \triangleright g \mapsto T^{\text{hor}} g = dL_{g}(\mathfrak{g}_{p^{(k)}}^\perp)$ defines a connection on the principal fiber bundle $G \to G / G_{p^{(k)}}$ for all $p \in M_k(H_k)$. Consider next the isotropy bundle over $M_k(H_k)$

$$\text{Iso} M_k(H_k) \to M_k(H_k),$$

as well as the canonical projection

$$\pi_k : W_k \to M_k(H_k), \quad f_k(p^{(k)}, v^{(k)}) \mapsto p^{(k)}, \quad p^{(k)} \in M_k(H_k), v^{(k)} \in (\nu_k)_{p^{(k)}}.$$

Since $g \in G$ is an isometry, the theorem of Whitehead implies

$$f_k(p^{(k)}, v^{(k)}) = g \cdot f_k(p^{(k)}, v^{(k)}) = (\exp_{g \cdot p^{(k)}} \circ \gamma_k)(g_{*, p^{(k)}}^{*}(v^{(k)})) \quad \Leftrightarrow \quad p^{(k)} = gp^{(k)}, v^{(k)} = g_{*, p^{(k)}}^{*}(v^{(k)}),$$

so that one concludes

$$\mathcal{N} \subset \text{Iso} W_L \cup \bigcup_{k=1}^{L-1} \pi_k^{*} \text{Iso} M_k(H_k),$$

where $\text{Iso} W_L \to W_L$ is the isotropy bundle over $W_L$, and

$$\pi_k^{*} \text{Iso} M_k(H_k) = \left\{ (f_k(p^{(k)}, v^{(k)}), h^{(k)}) \in W_k \times G : h^{(k)} \in G_{p^{(k)}} \right\}$$

denotes the induced bundle. Consider now an integral $I(\mu)$ of the form [3]. Introduce a partition of unity $\{ \chi_k \}_{k=1,\ldots,l}$, subordinated to the covering $M = W_1 \cup \cdots \cup W_L$, and define

$$I_k(\mu) = \int_{T^*Y} \int_{G} e^{i\mu \Phi(x, \xi, g)} \chi_k(x) a(gx, x, \xi, g) \, dg \, d(T^*Y)(x, \xi).$$

As will be explained in Lemma [3], the critical set of the phase function $\Phi$ is clean on the support of $\chi_{La}$, so that one can directly apply the stationary phase principle to obtain asymptotics for $I_L(\mu)$. 
We shall therefore turn to the case when $1 \leq k \leq L - 1$, and $W_k \cap Y \neq \emptyset$. Let $\{v_1^{(k)}, \ldots, v_{c}^{(k)}\}$ be an orthonormal frame in $\nu_k$, $(p_1^{(k)}, \ldots, p_{n-c}^{(k)})$ be local coordinates on $M_k(H_k)$, and write

$$\gamma^{(k)}(v^{(k)})(p^{(k)}), \theta^{(k)}) = \sum_{i=1}^{c} \theta_i^{(k)} v_i^{(k)}(p^{(k)}) \in \gamma^{(k)}(\nu_k)_{p^{(k)}}.$$  

(23)

By choosing $Y$ small enough, we can assume that the coordinates in the chart $(\kappa, f_k(\nu_k) \cap Y)$ are given by $\kappa(\exp_{p^{(k)}} \gamma^{(k)}(v^{(k)})) = (x_1, \ldots, x_n) = (p_1^{(k)}, \ldots, p_{n-c}^{(k)}, \theta_1^{(k)}, \ldots, \theta_{c}^{(k)})$. By the considerations leading to (22),

$$\text{Crit}_k(\Phi) \subset \pi^*\text{Iso}M_k(H_k) \times \mathbb{R}^n,$$

where

$$\text{Crit}_k(\Phi) = \{(x, \xi, g) \in (\Omega \cap T^*(W_k \cap Y)) \times G : g \cdot (x, \xi) = (x, \xi)\}.$$

Let therefore $U_k$ be a tubular neighborhood of $\pi^*\text{Iso}M_k(H_k)$ in $W_k \times G$, and

$$\Pi_k : U_k \rightarrow \pi_k^*\text{Iso}M_k(H_k)$$

the canonical projection which is obtained by considering geodesic normal coordinates around $\pi_k^*\text{Iso}M_k(H_k)$ and by identifying $U_k$ with a neighborhood of the zero section in the normal bundle $N \pi_k^*\text{Iso}M_k(H_k)$. The non-stationary phase theorem [28], Theorem 7.7.1, then yields

$$I_k(\mu) = \int_{U_k} \int_{\mathbb{R}^n} e^{i\Phi(x, \xi, g)} \chi_k(x) b(g, x, \xi, g) \, d\xi \, dg \, dM(x) + O(\mu^{-\infty}),$$

(24)

where $b$ is equal to the amplitude $a$ multiplied by a smooth cut-off-function with compact support in $U_k$. Note that the fiber of $N \pi^*\text{Iso}M_k(H_k)$ at a point $(f_k(p^{(k)}), v^{(k)}), h^{(k)})$ may be identified with the fiber of the normal bundle to $G_{p^{(k)}}$ at the point $h^{(k)}$. Let now $A_1(p^{(k)}), \ldots, A_d(p^{(k)})$ be an orthonormal basis of $g_{p^{(k)}}$, and $B_1(p^{(k)}), \ldots, B_{c}^{(k)}(p^{(k)})$ an orthonormal basis of $g_{p^{(k)}}$, and introduce canonical coordinates of the second kind

$$\left(\alpha_1, \ldots, \alpha_d^{(k)}, \beta_1, \ldots, \beta_{c}^{(k)}\right) \mapsto \sum_{i=1}^d \alpha_i A_i(p^{(k)}) \sum_{j=1}^c \beta_j B_j(p^{(k)}) \, g,$$

in a neighborhood of a point $g \in G$, see [25], page 146, which in turn give rise to coordinates

$$\left(\alpha_1, \ldots, \alpha_d^{(k)}\right) \mapsto (f_k(p^{(k)}), v^{(k)}), \sum_{i=1}^d \alpha_i A_i(p^{(k)}) \, h^{(k)}),$$

in $\Pi_k^{-1}(f_k(p^{(k)}), v^{(k)}), h^{(k)})$. Integrating along the fibers of the normal bundle to $\pi_k^*\text{Iso}M_k(H_k)$, compare [14], page 30, we obtain for $I_k(\mu)$ the expression

$$I_k(\mu) = \int_{\pi_k^*\text{Iso}M_k(H_k)} \int_{\pi_k^{-1}(f_k(p^{(k)}, v^{(k)}), h^{(k)}) \times \mathbb{R}^n} e^{i\Phi} \chi_k b \, J_k \, d\xi \, dA^{(k)}) \, dh^{(k)} \, dv^{(k)} \, dp^{(k)}$$

(25)

$$= \int_{M_k(H_k)} \int_{\pi_k^{-1}(p^{(k)}) \times G_{p^{(k)}} \times D_{1}\left(g_{p^{(k)}}^{1}\right) \times \mathbb{R}^n} e^{i\Phi} \chi_k b \, J_k \, d\xi \, dA^{(k)} \, dh^{(k)} \, dv^{(k)} \, dp^{(k)}$$

up to a term of order $O(\mu^{-\infty})$, where $dp^{(k)}, dv^{(k)}, dh^{(k)}, dA^{(k)}$ are suitable volume densities on the sets $M_k(H_k), \nu_k, G_{p^{(k)}}, g_{p^{(k)}}^{1} \simeq N_{h^{(k)}}G_{p^{(k)}}$, respectively, and

$$\left(p^{(k)}(k), v^{(k)}, A^{(k)}, h^{(k)}\right) \mapsto (f_k(p^{(k)}), v^{(k)}), e^{A^{(k)}} h^{(k)} = (x, g),$$

are coordinates on $U_k$ such that $dg \, dM(x) = J_k \, dA^{(k)} \, dh^{(k)} \, dv^{(k)} \, dp^{(k)}$, $J_k$ being a Jacobian. Here $D_{1}\left(g_{p^{(k)}}^{1}\right)$ denotes the interior of a ball of suitable radius $\epsilon > 0$ around the origin in $g_{p^{(k)}}^{1}$.
First monoidal transformation. We shall now successively resolve the singularities of \((12)\). To begin with, note that (22) implies
\[
\mathcal{N} = \text{Iso} W_L \cup \bigcup_{k=1}^{L-1} \mathcal{N} \cap U_k,
\]
and we put \(\mathcal{N}_L = \text{Iso} W_L, \mathcal{N}_k = \mathcal{N} \cap U_k\). While \(\mathcal{N}_L\) is a smooth submanifold, \(\mathcal{N}_k\) is in general singular. In particular, if \(\dim H_k \neq \dim H_L, \mathcal{N}_k\) has a singular locus given by \(\text{Iso} M_k(H_k)\). We shall therefore perform for each \(k \in \{1, \ldots, L - 1\}\) a monoidal transformation
\[
\zeta_k : B_{Z_k}(U_k) \longrightarrow U_k
\]
with center \(Z_k = \text{Iso} M_k(H_k) \subset \mathcal{N}_k\). By piecing these transformations together, we obtain the monoidal transformation
\[
\zeta^{(1)} : B_{Z^{(1)}}(\mathcal{M}) \longrightarrow \mathcal{M}, \quad Z^{(1)} = \bigcup_{k=1}^{L-1} Z_k \quad \text{(disjoint union)}.
\]
To get a local description, let \(k\) be fixed, and write \(A^{(k)}(p^{(k)}, \alpha^{(k)}) = \sum \alpha_i^{(k)} A_i^{(k)}(p^{(k)}) \in \mathfrak{g}^+_p(\mathfrak{q})\), \(B^{(k)}(p^{(k)}, \beta^{(k)}) = \sum \beta_j^{(k)} B_j^{(k)}(p^{(k)}) \in \mathfrak{g}^p(\mathfrak{q})\). With respect to these coordinates and the ones introduced in (23) and (20) we have \(Z_k \simeq \{ T^{(k)} = (\theta^{(k)}, \alpha^{(k)}) = 0 \}\), so that
\[
B_{Z_k}(U_k) = \left\{ (x, g, [t]) \in U_k \times \mathbb{R}P^{(k)+d(k)-1} : T_i^{(k)} t_j^{(k)} = T_j^{(k)} t_i^{(k)} \right\},
\]
\[
\zeta_k : (x, g, [t]) \mapsto (x, g).
\]
If \(t_0 \neq 0\),
\[
(x, g, [t]) \mapsto \left( p^{(k)}, h^{(k)}, \frac{t_1}{t_0}, \ldots, \frac{t_{c(k)+d(k)} - 1}{t_0}, T^{(k)} \right)
\]
define local coordinates on \(B_{Z_k}(U_k)\). Consequently, setting \(V_0 = \{ [t] \in \mathbb{R}P^{(k)+d(k)-1} : t_0 \neq 0 \}\), we can cover \(B_{Z_k}(U_k)\) with charts \(\{(\varphi^0_k, \mathcal{O}^0_k)\}\), where \(\mathcal{O}^0_k = B_{Z_k}(U_k) \cap (U_k \times V_0)\), such that \(\zeta_k\) is realized in each of the \(\theta^{(k)}\) charts \(\mathcal{O}^0_k\) as
\[
\zeta_{0} = \zeta^{(1)} \circ \varphi^{0}_{k} = 1 : (p^{(k)}, r_k, \tilde{v}^{(k)}, A^{(k)}, h^{(k)}) \mapsto (p^{(k)}, r_k \tilde{v}^{(k)}, \tau_k A^{(k)}, h^{(k)})
\]
\[
\mapsto (\exp_{p^{(k)}} r_k \tilde{v}^{(k)}, e^{\tau_k A^{(k)}} h^{(k)}) = (x, g),
\]
where \(\tilde{v}^{(k)}(p^{(k)}, \theta^{(k)}) \in \gamma^{(k)}((S^+_k)_p(\mathfrak{q}))\), and \(S^+_k = \left\{ v \in \nu_k : v = \sum s_i v_i^{(k)}, s > 0, \|v\| = 1 \right\}\), while \(\tau_k \in (-1, 1)\). Note that for each \(1 \leq \varrho \leq c(k)\) we have \(W_k \simeq S^+_k \times (-1, 1)\) up to a set of measure zero. A similar description of \(\zeta^{(1)}\) is given in the \(\alpha^{(k)}\) charts. As a consequence, we obtain a partial monomialization of the inverse image ideal sheaf \((\zeta^{(1)})^*(I_\mathcal{N})\)
\[
(\zeta^{(1)})^*(I_\mathcal{N}) \cdot \mathcal{E}_{(x, g, [t]), B_{Z^{(1)}}(\mathcal{M})} = \tau_k : (\zeta^{(1)})^{-1}(I_\mathcal{N}) \cdot \mathcal{E}_{(x, g, [t]), B_{Z^{(1)}}(\mathcal{M})}
\]
in a neighborhood of any point \((x, g, [t]) \in B_{Z^{(1)}}(\mathcal{M})\). To see this, note that \(I_\mathcal{N}\) is generated locally by the functions \(\tilde{x}_q(x) - \tilde{x}_q(g \cdot x), 1 \leq q \leq n\). We have \(g \cdot \exp_{p^{(k)}} r_k \tilde{v}^{(k)} = \exp_{p^{(k)}} [g_{*p^{(k)}}(r_k \tilde{v}^{(k)})]\), where \(g_{*p^{(k)}}(\tilde{v}^{(k)}) \in \gamma^{(k)}((\nu_k)_{p^{(k)}})\), \(\nu_k\) being a G-vector bundle. Now, Taylor expansion at \(\tau_k = 0\) gives for \(y \in Y \cap f_k(\nu_k)\)
\[
\tilde{x}_q(e^{\tau_k A^{(k)}} \cdot y) = \tilde{x}_q(y) - r_k \tilde{A}_y^{(k)}(\tilde{x}_q) + O(|\tau_k^2 A^{(k)}|),
\]
where $\tau_k \in (-1, 1)$, $A^{(k)} \in D_\nu (g_{\mu^{(k)}}^\nu)$, and $\iota > 0$ is assumed to be sufficiently small. Furthermore, $A^{(k)}_y (\tilde{x}_q) = d\tilde{x}_q (A^{(k)}_y)$. Consequently,

$$
\kappa \left( \exp_{\nu^{(k)}} \tau_k \tilde{v}^{(k)} \right) - \kappa \left( e^{\tau_k A^{(k)}} h^{(k)} \cdot \exp_{\nu^{(k)}} \tau_k \tilde{v}^{(k)} \right) = \tau_k \left( A^{(k)}_y (\tilde{x}_1), \ldots, \theta_1^{(k)} \left( \tilde{v}^{(k)} \right) - \theta_1^{(k)} \left( (h^{(k)})_{\nu^{(k)}} \tilde{v}^{(k)} \right), \ldots \right) + O(\left| \tau_k A^{(k)} \right|).
$$

(28)

Since similar considerations hold in the $\alpha^{(k)}$-charts $\{O^{(k)}_k\}_{c^{(k)} + 1 \leq \theta \leq c^{(k)} + d^{(k)}}$, the assertion follows. In the same way, the phase function (10) factorizes according to

$$
\Phi \circ (\text{id} \otimes \zeta^0_k) = (k) \tilde{\Phi}_{\text{tot}} = \tau_k \cdot (k) \tilde{\Phi}_{\text{wk}},
$$

(29)$\tilde{\Phi}_{\text{tot}}$ and $\tilde{\Phi}_{\text{wk}}^{(k)}$ being the total and weak transform of the phase function $\Phi$, respectively. In the $\theta^{(k)}$-charts this explicitly reads

$$
\Phi (x, \xi, g) = \left( \kappa \left( \exp_{\nu^{(k)}} \tau_k \tilde{v}^{(k)} \right) - \kappa \left( e^{\tau_k A^{(k)}} h^{(k)} \cdot \exp_{\nu^{(k)}} \tau_k \tilde{v}^{(k)} \right), \xi \right) = \tau_k \left[ \sum_{q=1}^{n-c^{(k)}} \xi_q dp_q^{(k)} (A^{(k)}_y (\nu^{(k)})) + \sum_{r=1}^{c^{(k)}} \left( \theta_r^{(k)} (\tilde{v}^{(k)}) - \theta_r^{(k)} ((h^{(k)})_{\nu^{(k)}} \tilde{v}^{(k)}) \right) \xi_{n-c^{(k)}+r} + O(\left| \tau_k A^{(k)} \right|) \right].
$$

(30)

Since $\zeta_k$ is a real analytic, surjective map, we can lift the integral $I_k (\mu)$ to the resolution space $B_{\zeta_k} (U_k)$, and introducing a partition $\{u^{(k)}_\mu\}$ of unity subordinate to the covering $\{O^{(k)}_k\}$ yields with (24) the equality

$$
I_k (\mu) = \sum_{\varphi^{(k)}_\mu} I^{(k)}_\mu + \sum_{\varphi^{(k)}_\mu = c^{(k)} + 1} \tilde{I}^{(k)}_\mu (\mu)
$$

up to terms of order $O(\mu^{-\infty})$, where the integrals $I^{(k)}_\mu (\mu)$ and $\tilde{I}^{(k)}_\mu (\mu)$ are given by the expressions

$$
\int_{B_{\zeta_k} (U_k) \times \mathbb{R}^n} u^{(k)}_\mu (\text{id} \otimes \zeta_k)^* (e^{i\mu \Phi} \chi_k b) \, dg \, dM (x) \, d\xi.
$$

As we shall see, the weak transforms $(k) \tilde{\Phi}_{\text{wk}}^{(k)}$ have no critical points in the $\alpha^{(k)}$-charts, which will imply that the integrals $I^{(k)}_\mu (\mu)$ contribute to $I (\mu)$ only with lower order terms. In what follows, we shall therefore restrict ourselves to the examination of the integrals $I^{(k)}_\mu (\mu)$. Setting $\alpha^{(k)}_\mu = (u_\varphi \circ (\varphi^{(k)}_\mu)^{-1}) \cdot [(b \chi_k) \circ (\text{id} \otimes \zeta_k)]$ we obtain with (25) and (27)

$$
I^{(k)}_\mu (\mu) = \int_{M_k (H_k) \times (-1, 1)} \left[ \int_{\gamma^{(k)} ((S^k)_{\mu^{(k)}}) \times G_{\mu^{(k)}} \times D_{\nu^{(k)}} (g_{\mu^{(k)}}^\nu) \times \mathbb{R}^n} e^{i\mu \tau_k (k) \tilde{\Phi}_{\text{wk}}^{(k)}} u^{(k)}_\mu \tilde{J}_k \right] d\xi \, dA^{(k)} \, dh^{(k)} \, d\tilde{v}^{(k)} \, d\tau_k \, dp^{(k)},
$$

where $d\tilde{v}^{(k)}$ is a suitable volume density on $\gamma^{(k)} ((S^k)_{\mu^{(k)}})$ such that the pulled back density reads $(\zeta_k^0)^* (dg \, dM (x)) = \tilde{J}_k^0 dA^{(k)} dh^{(k)} d\tilde{v}^{(k)} d\tau_k dp^{(k)}$. Furthermore, by comparing (20) and (27) one sees that

$$
\tilde{J}_k^0 = |\tau_k|^{c^{(k)} + d^{(k)} - 1} J_k \circ \zeta_k^0.
$$

Note that the weak transform is defined only locally, while the total transform has a global meaning. To keep the notation as simple as possible, we restrained ourselves from making the chart dependence of $\tau_k$ and $(k) \tilde{\Phi}_{\text{wk}}^{(k)}$ manifest.
First reduction. Let $k$ be fixed, and assume that there exists a $x \in W_k$ with isotropy group $G_x \sim H_j$, and let $p^{(k)} \in M_k(H_k), v^{(k)} \in (\nu_k)_p^{(k)}$ be such that $x = f_k(p^{(k)}, v^{(k)})$. Since we can assume that $x$ lies in a slice at $p^{(k)}$ around the $G$-orbit of $p^{(k)}$, we have $G_x \subset G_{p^{(k)}}$, see [30], pp. 184, and [31], page 86. Hence $H_j$ must be conjugate to a subgroup of $H_k \sim G_{p^{(k)}}$. Now, $G$ acts on $M_k$ with the isotropy types $(H_k, (H_{k+1}), \ldots, (H_L))$. The isotropy types occurring in $W_k$ are therefore those for which the corresponding isotropy groups $H_k, H_{k+1}, \ldots, H_L$ are conjugate to a subgroup of $H_k$, and we shall denote them by $(H_k) = (H_{l_1}), (H_{l_2}), \ldots, (H_L)$. By the invariant tubular neighborhood theorem, one has the isomorphism

$$W_k/G \cong (\nu_k)^{p^{(k)}}_v / G_{p^{(k)}}$$

for every $p^{(k)} \in M_k(H_k)$. Furthermore, $(\nu_k)^{p^{(k)}}_v$ is an orthogonal $G_{p^{(k)}}$-space; therefore $G_{p^{(k)}}$ acts on $(S_k)^{p^{(k)}}$ with isotropy types $(H_{l_1}), \ldots, (H_L)$, cp. [14], pp. 34, and $G$ must act on $S_k$ with isotropy types $(H_{l_1}), \ldots, (H_L)$ as well. If all isotropy groups $H_{l_1}, \ldots, H_L$ have the same dimension, the singularities of $N_k$ have been resolved. Indeed, note that $\zeta_k^{-1}(N_k)$ is contained in the union of the $\theta^{(k)}$-charts $\{O_k^{(p)}\}_{1 \leq p \leq e^{(k)}}$ since, in the notation of [20], $e^{(k)}h^{(k)} \in G_{f_k(p^{(k)}, v^{(k)})} \subset G_{p^{(k)}}$ necessarily implies $A^{(k)} = 0$. Let therefore $1 \leq \rho \leq e^{(k)}$, and consider the set $\zeta_k^{-1}(N_k) \cap O_k^{\rho}$, which is given by all points $(x, g, [\ell])$ with coordinates $(p^{(k)}, \tau_k, \nu^{(k)}, A^{(k)}, h^{(k)})$ satisfying

$$e^{\tau_k}A^{(k)}h^{(k)} \in G_{\exp_{p^{(k)}}(\nu^{(k)}, [\ell])} \subset G_{p^{(k)}}.$$ 

If $\tau_k \neq 0$, this implies $A^{(k)} = 0$ and $h^{(k)} \in G_{\mu^{(k)}}$. Therefore

$$\zeta_k^{-1}(N_k) \cap O_k^{\rho} = \{A^{(k)} = 0, h^{(k)} \in G_{\mu^{(k)}}, \tau_k \neq 0 \} \cup \{\tau_k = 0\}.$$ 

Assume now that all isotropy groups $H_{l_1}, \ldots, H_L$ have the same dimension. If $H_k$ has the same dimension, too, $N_k$ is already a manifold. Otherwise, the invariant tubular neighborhood theorem implies that $\zeta_k^{-1}(\text{Reg} \cap N_k) \cap O_k^{\rho} = \{A^{(k)} = 0, h^{(k)} \in G_{\mu^{(k)}}, \tau_k \neq 0 \}$, where $\text{Reg} \cap N_k = \text{Reg} \cap U_k$ denotes the regular part of $N_k$. The closure of this set is a smooth manifold, and taking the union over all $1 \leq \rho \leq e^{(k)}$ yields a smooth manifold $\tilde{N}_k \subset BZ_k(U_k)$ which intersects $\zeta_k^{-1}(\text{Sing} \cap N_k)$ normally. After performing an additional monomial transformation with center $\tilde{N}_k \cap \zeta_k^{-1}(\text{Sing} \cap N_k)$, we obtain a strong resolution for $\tilde{N}_k$. Furthermore, if $G$ acts on $S_k$ only with isotropy type $(H_L)$, we shall see in Sections [6] and [7] that in each of the $\theta^{(k)}$-charts the critical sets of the weak transforms $(k)\tilde{\phi}^{wk}$ are clean, so that one can apply the stationary phase theorem in order to compute each of the $I^p_k(\mu)$. But in general, $G$ will act on $S_k$ with singular orbit types, so that neither $\tilde{N}_k$ is resolved, nor do the weak transforms $(k)\tilde{\phi}^{wk}$ have clean critical sets, and we are forced to continue with the iteration.

Second Decomposition. In what follows, let $1 \leq k \leq L - 2$, and $p^{(k)} \in M_k(H_k)$ be fixed. Since $\gamma^{(k)} : \nu_k \to \nu_k$ is an equivariant diffeomorphism onto its image, $\gamma^{(k)}(S_k)^{p^{(k)}}$ is a compact $G_{p^{(k)}}$-manifold, and we consider the covering

$$\gamma^{(k)}(S_k)^{p^{(k)}} = W_{kl_1} \cup \cdots \cup W_{kl_k}, \quad W_{kl_j} = f_{kl_j}(\hat{D}_1(\nu_{kl_j})), \quad W_{kl_L} = \text{Int}(\gamma^{(k)}(S_k)^{p^{(k)}}L_k),$$

where $f_{kl_j} : \nu_{kl_j} \to \gamma^{(k)}(S_k)^{p^{(k)})}l_j$ is an invariant tubular neighborhood of $\gamma^{(k)}(S_k)^{p^{(k)}}l_j(H_{l_j})$ in

$$\gamma^{(k)}(S_k)^{p^{(k)}}l_j = \gamma^{(k)}(S_k)^{p^{(k)} - j} - \bigcup_{r=2}^{j-1} f_{kl_r}(\hat{D}_{l_2/2}(\nu_{kl_r})), \quad j \geq 2,$$
and $f_{kl}(p^{(l)}_i, v^{(l)}_i) = (\exp_{v^{(l)}_i} \circ \gamma^{(l)}_i)(v^{(l)}_i), p^{(l)}_i \in \gamma^{(l)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}), v^{(l)}_i \in (\nu_{klj})_{p^{(ij)}_j}, \gamma^{(l)}_i : \nu_{klj} \to \nu_{klj},$ being an equivariant diffeomorphism onto its image given by

$$
\gamma^{(l)}_i(v^{(l)}_i) = \frac{F_{lj}(p^{(lj)}_i)}{(1 + \|v^{(l)}_i\|)^1/2} v^{(l)}_i,
$$

where $F_{lj} : ((S_k)_{p^{(k)}})_{l_j}(H_{l_j}) \to \mathbb{R}$ is a smooth, $G_{p^{(k)}}$-invariant, positive function. Let now $\{\chi_{klj}\}$ denote a partition of unity subordinate to the covering $\{W_{klj}\}$, which extends to a partition of unity on $\gamma^{(k)}(S_k)$ as a consequence of the invariant tubular neighborhood theorem, by which in particular $\gamma^{(k)}(S_k)/G \simeq \gamma^{(k)}((S_k)_{p^{(k)}})/G_{p^{(k)}}$ for all $p^{(k)}$. We then define

$$I^0_{klj}(\mu) = \int_{M_k(H_k) \times (-1,1)} \int_{\gamma^{(k)}((S_k)_{p^{(k)}}) \times G_{p^{(k)}} \times D_1(\nu_{klj}) \times \mathbb{R}^n} e^{i\mu \tau_k^{(l)} \Phi^{(l)}} d\mu \cdot d\Phi^{(k)} d\tau_k d\nu_{klj},$$

so that $I^0_{klj}(\mu) = I^0_{klj}(\mu) + \cdots + I^0_{klj}(\mu)$. Since $G_{p^{(k)}}$ acts on $W_{klj}$ only with type $(H_L)$, the iteration process for $I^0_{klj}(\mu)$ ends here. For the remaining integrals $I^0_{klj}(\mu)$ with $k < l < L$ and non-zero integrand, let us denote by

$$\text{Iso} \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}) \to \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j})$$

the isotropy bundle over $\gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j})$, and by $\pi_{klj} : W_{klj} \to \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j})$ the canonical projection. We then assert that in each $\theta^{(k)}$-chart \{\mathcal{O}_{klj}^{(k)}\}_{1 \leq \rho \leq c^{(k)}}

$$\text{Crit}_{klj}^{(k)}(\Phi^{(l)}) \
\subset \left\{ (p^{(k)}, \tau_k, \Phi^{(l)}_k, \xi, h^{(k)}, A^{(k)} : (\Phi^{(l)}_k, h^{(k)}) \in \pi_{klj}^{*}\text{Iso} \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}), \text{A}^{(k)} = 0 \right\},$$

where

$$\text{Crit}_{klj}^{(k)}(\Phi^{(l)}) = \left\{ (p^{(k)}, \tau_k, \Phi^{(l)}_k, \xi, h^{(k)}, A^{(k)} : (\Phi^{(l)}_k, h^{(k)}) = 0, \Phi^{(l)}_k \in W_{klj} \right\}.$$ 

Indeed, from (29) it is clear that for $\tau_k \neq 0$ the condition $\partial^{(k)}(\Phi^{(l)}) = 0$ is equivalent to

$$e^{\tau_k \sum \alpha^{(k)} A^{(k)}(p^{(k)}, h^{(k)})} \in G_{\exp_{p^{(k)}} \tau_k \phi^{(k)}} \subset G_{p^{(k)}},$$

which implies $\alpha^{(k)} = 0$, and consequently $h^{(k)} \in G_{\phi^{(k)}}$. But if $\tilde{\Phi^{(k)}} = f_{klj}(p^{(l)}_i, v)$, where $p^{(l)}_i \in \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}), v \in D_1(\nu_{klj})_{p^{(ij)}_j}$, then $h^{(k)} \in G_{p^{(l)}}$. On the other hand, assume that $\tau_k = 0$. By (29), the vanishing of the $\zeta$-derivatives of $(\Phi^{(l)}) = 0,$ which again implies $\alpha^{(k)} = 0$, as well as $h^{(k)} \in G_{\phi^{(k)}}$. But if $\tilde{\Phi^{(k)}} = f_{klj}(p^{(l)}_i, v)$ as above, we again conclude $h^{(k)} \in G_{p^{(l)}}$, and (32) follows, since

$$\pi_{klj}^{*}\text{Iso} \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}) = \{ (w, g) \in W_{klj} \times G_{p^{(k)}} : w = f_{klj}(p^{(l)}_i, v), p^{(l)}_i \in \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}), v \in D_1(\nu_{klj})_{p^{(ij)}_j}, g \in G_{p^{(l)}} \}.$$ 

The same reasoning also shows that the weak transforms $(\Phi^{(l)}) = 0,$ can have no critical points in the $\alpha^{(k)}$-charts \{\mathcal{O}_{klj}^{(k)}\}_{1 \leq \rho \leq c^{(k)} + d^{(k)}}. \text{ Let now U}_{klj} denote a tubular neighborhood of the set } \pi_{klj}^{*}\text{Iso} \gamma^{(k)}((S_k)_{p^{(k)}})_{l_j}(H_{l_j}) \text{ in } W_{klj} \times G_{p^{(k)}}, \text{ and let } b^{(k)}_l = \text{equal to the product of the amplitude } a^{(k)}_l \text{ with some smooth cut-off-function with compact support in } U_{klj} \text{ that depends smoothly on } p^{(k)}. \text{ The non-stationary phase theorem then implies that, up to terms of lower order, we can replace }
If we now cover \( \theta \), consider next the decomposition
\[
\mathfrak{g} = \mathfrak{g}_{p(k)} \oplus \mathfrak{g}_{p(k)}^\perp = (\mathfrak{g}_{p(i)} \oplus \mathfrak{g}_{p(i)}^\perp) \oplus \mathfrak{g}_{p(i)}^\perp.
\]
Let further \( h^{(l)} \in G_{p(l)} \), and \( A_1^{(l)}, \ldots, A_{d_p(l)}^{(l)} \) be an orthonormal frame in \( \mathfrak{g}_{p(i)}^\perp \), as well as \( B_1^{(l)}, \ldots, B_{d_p(l)}^{(l)} \) be an orthonormal frame in \( \mathfrak{g}_{p(i)} \), and \( v_{i_1}^{(k)}, \ldots, v_{i_{d_p(l)}}^{(k)} \) an orthonormal frame in \( (\nu_{kl})_{p(i)} \). Integrating along the fibers in a neighborhood of \( \pi_{klj}^* \text{Iso} \gamma^{(k)}(\mathcal{S}_{k(p(k)})_{l_j}(H_i)) \) then yields for \( \mathcal{I}_{klj}^p(\mu) \) the expression
\[
\mathcal{I}_{klj}^p(\mu) = \int_{\mathcal{M}_k(H_L) \times (\nu_{klj})_{p(l)}} \left[ \int_{\mathcal{M}_{klj}(H_i)} \left[ \int_{\mathcal{F}_{klj}^p} \left[ \frac{e^{i\mu\tau_k^{(k)}(H)}\mathfrak{g}_{w(k)}}{p(l)} \right] \right] \right] d\tau_k d\nu_{klj}
\]
up to lower order terms, where \( \mathcal{F}_{klj}^p \) is a Jacobian, and
\[
(p^{(l)}, v^{(l)}, A^{(l)}, h^{(l)}) \mapsto \left( f_{klj}(p^{(l)}, v^{(l)}), e^{A^{(l)}}, h^{(l)} \right) = \left( \tilde{\varphi}^{(k)}, \tilde{h}^{(k)} \right)
\]
are coordinates on \( U_{klj} \), while \( dp^{(l)}, dA^{(l)}, dh^{(l)} \), and \( dv^{(l)} \) are suitable volume densities in the spaces \( \gamma^{(k)}(\mathcal{S}_{k(p(k)})_{l_j}(H_i)), \mathfrak{g}_{p(i)}^\perp, G_{p(i)}, \) and \( \mathcal{D}_1(\nu_{klj})_p(l) \), respectively, such that we have the equality \( \mathcal{F}_{klj}^p dh^{(k)} d\tilde{v}^{(k)} \equiv \mathcal{F}_{klj}^p dA^{(l)} dh^{(l)} dv^{(l)} dp^{(l)} \).

**Second monoidal transformation.** Put \( \tilde{M}^{(1)} = B_{2^{(1)}}(\mathcal{M}) \), and consider the monoidal transformation
\[
\zeta^{(2)} : B_{2^{(2)}}(\tilde{M}^{(1)}) \rightarrow \tilde{M}^{(1)}, \quad Z^{(2)} = \bigcup_{k<l \leq L \cap (H_i) \subseteq (H_k)} Z_{kl}, \quad (\text{disjoint union}),
\]
where
\[
Z_{kl} \simeq \bigcup_{p^{(k)} \in M_{klj}(H_k)} (\nu_{klj})_{p(l)}, \quad k < l \leq L, \quad (H_i) \subseteq (H_k),
\]
are the possible maximal singular loci of \( \zeta^{(1)}_k^{-1}(\mathcal{N}) \). To obtain a local description of \( \zeta^{(2)} \), let us write \( A^{(l)}(p^{(k)}, p^{(l)}, \alpha^{(l)}) = \alpha_1^{(l)} A_1^{(l)}(p^{(k)}, p^{(l)}) \in \mathfrak{g}_{p^{(l)}}^\perp, B^{(l)}(p^{(k)}, p^{(l)}, \beta^{(l)}) = \sum_{i=1}^{d_p(l)} B_{i}^{(l)}(p^{(k)}, p^{(l)}) \in \mathfrak{g}_{p^{(l)}} \), as well as
\[
\gamma^{(l)}(v^{(l)})(p^{(k)}, p^{(l)}, \theta^{(l)}) = \sum_{i=1}^{d_p(l)} \theta_i^{(l)} v_i^{(kl)}(p^{(k)}, p^{(l)}) \in \gamma^{(l)}(\nu_{klj})_{p^{(l)}}.
\]
One has \( Z_{kl} \simeq \{ \alpha^{(k)} = 0, \alpha^{(l)} = 0, \theta^{(l)} = 0 \} \), which in particular shows that each \( Z_{kl} \) is a manifold.

If we now cover \( B_{2^{(2)}}(\tilde{M}^{(1)}) \) with the standard charts, a computation shows that \( (\zeta^{(1)} \circ \zeta^{(2)})^{-1}(\mathcal{N}) \) is contained in the \( \{ \theta^{(k)}, \theta^{(l)} \} \)-charts. For our purposes, it will therefore suffice to examine \( \zeta^{(2)} \) in each of these charts in which it reads
\[
\zeta^{(2)}_{kl} : (p^{(k)}, \tau_k, p^{(l)}, \tau_l, v^{(l)}, A^{(l)}, h^{(l)}, A^{(k)}) \mapsto \begin{cases} 
\gamma^{(l)}(p^{(k)}, \tau_k, p^{(l)}, \tau_l, v^{(l)}, A^{(l)}, h^{(l)}, A^{(k)}) 
\end{cases}
\]
\[
\tau_k, \tau_l, v^{(l)}, e^{A^{(l)}}, h^{(l)}, A^{(k)} \end{cases} \end{cases}
\]
where \( \tau \in (-1, 1) \), and
\[
\tilde{v}^{(l)}(p^{(k)}, p^{(l)}, \theta^{(l)}) \in \gamma^{(l)}((\mathcal{S}_{klj}^+)_{p^{(l)}}).
Here $S_{kl}$ stands for the the sphere subbundle in $\nu_{kl}$, and $S^+_{kl} = \{ v \in S_{kl} : v = \sum s_i v_i^{(kl)}, v_\sigma > 0 \}$ for some $\sigma$. Note that $Z_{kl}$ has normal crossings with the exceptional divisor $E_k = \zeta_k^{-1}(Z_k) \simeq \{ t_k = 0 \}$, and that for each $p^{(k)} \in M_k(H_k)$ we have $W_k \simeq S^+_kl \times (-1, 1)$, up to a set of measure zero. Now, Taylor expansion at $\tau_l = 0$ gives

$$T \theta^{(k)} \left( \exp_{p^{(k)}} (\tau_l \theta^{(l)}) - T \theta^{(k)} \left( (e^{\tau_l A^{(l)}} h^{(l)})_{*p^{(k)}} \exp_{p^{(k)}} (\tau_l \theta^{(l)}) \right) \right) = \tau_l \left( \frac{\partial \theta^{(k)}(1, p^{(l)}, 0)}{\partial (\tau_k, p^{(l)}, \theta^{(l)})} \right) \tau_l \left( 0, d\tau_1^{(l)}, \ldots, d\tau_{e(l)_{*c(l)-c(l)-1}}^{(l)}(A^{(l)}_{p^{(l)}}, \theta^{(l)}(\tilde{\theta^{(l)}) \theta^{(l)}((h^{(l)})_{*p^{(k)}} \tilde{\theta^{(l)})}, \ldots, \theta^{(l)}(\tilde{\theta^{(l)}) \theta^{(l)}((h^{(l)})_{*p^{(k)}} \tilde{\theta^{(l)})}, O(|\tau_l^2 A^{(l)}|) + O(|\tau_l^2 [\theta^{(l)}(\tilde{\theta^{(l)}) \theta^{(l)}((h^{(l)})_{*p^{(k)}} \tilde{\theta^{(l)})], \ldots, O(|\tau_l^2 A^{(l)}|) + O(|\tau_l^2 [\theta^{(l)}(\tilde{\theta^{(l)}) \theta^{(l)}((h^{(l)})_{*p^{(k)}} \tilde{\theta^{(l)})] \right) \right)$$

where $\{ p^{(l)} \}$ are local coordinates on $\gamma^{(k)}((S_{kl})_{p^{(k)}})(H_l)$.

Here $\tilde{\theta}$ denotes the Jacobian of the coordinate change $\phi^{(k)} = \theta^{(k)}(\tau_k \exp_{p^{(l)}} (\gamma^{(l)}(\mu^{(l)})))$, and all vectors are considered as row vectors, the transposed being a column vector. Since similar considerations hold in the other charts, we obtain with (28) and (33) a partial monomialization of $(\zeta^{(1)} \circ \zeta^{(2)})^{*}(I_N)$ according to 

$$(\zeta^{(1)} \circ \zeta^{(2)})^{*}(I_N) \cdot E_{m,B^*((M^{(l)})^1)} = \tau_k \tau_l \cdot (\zeta^{(1)} \circ \zeta^{(2)})^{*}(I_N) \cdot E_{m,B^*((M^{(l)})^1)}$$

in a neighborhood of any point $m \in B_{Z^{(2)}((M^{(l)})^1)}$. In the same way, the phase function factorizes locally according to 

$$\Phi \circ (\id \xi \otimes (\xi^{\rho} \circ \zeta^{\sigma})) = (kl) \tilde{\Phi}^{tot} = \tau_k \tau_l \tilde{\Phi}^{wk},$$

which by (30) and (33) explicitly reads

$$\Phi(x, \xi, g) = \tau_k \tau_l \sum_{q=1}^{c^{(k)}} \xi_q dp^{(k)}_q(A^{(l)}_{p^{(l)}})$$

$$+ \sum_{r=1}^{c^{(l)}} \left[ \theta^{(l)}_{r^{(l)}} \left( (e^{\tau_l A^{(l)}} h^{(l)})_{*p^{(k)}} \exp_{p^{(k)}} (\tau_l \theta^{(l)}) \right) \right] \xi_{n-c^{(l)}+r} + O(|\tau_l A^{(l)}|)$$

$$= \tau_k \tau_l \left[ \left( \frac{\partial(p^{(k)}, \theta^{(k)}(1, p^{(l)}, 0)}{\partial (p^{(k)}, \tau_k, p^{(l)}, \theta^{(l)})} \right) \left( dp^{(k)}(A^{(k)}_{p^{(k)}}, \ldots, dp^{(k)}(A^{(k)}_{p^{(k)}}, \ldots, dp^{(k)}(\tau_l \theta^{(l)}) - \theta^{(l)}(\tau_l \theta^{(l)}) \right) \right]$$

in the $(\theta^{(k)}, \theta^{(l)})$-charts. A computation now shows that the weak transforms $(kl) \tilde{\Phi}^{wk}$ have no critical points in the $(\theta^{(k)}, \alpha^{(l)})$-charts. We shall therefore see in Section 9 that modulo lower order terms $I^{\sigma}_{kl}(\mu)$ is given by a sum of integrals of the form

$$I^{\sigma}_{kl}(\mu) = \int_{M_k(H_k) \times (-1, 1)} \int_{\gamma^{(l)}((S_{kl})_{p^{(k)}})(H_l) \times (-1, 1)} \int_{\gamma^{(l)}((S_{kl})_{p^{(k)}} \times G_{p^{(k)}}(\tau_l \theta^{(l)}) \times \hat{D}_s(\sigma^{\sigma}_{p^{(k)}}) \times \hat{D}_s(\sigma^{\sigma}_{p^{(k)}}) \times \mathbb{R}^n}$$

$$e^{i\tau_l \tau_k \tilde{\theta}^{wk} a^{\sigma}_{kl} \tilde{\Sigma}^{\sigma}_{kl} d\xi dA^{(k)} dA^{(l)} d\tau_l d\tau_k dp^{(k)} dp^{(l)} d\tau_k dp^{(l)}$$







for some $t > 0$, where $a_{kl}^{(i)}$ are compactly supported amplitudes, and $d	ilde{\tau}^{(i)}$ is a suitable density on $\gamma^{(i)}((S_{kl})_{p^{(i)}})$ such that we have the equality

$$dM(x) dg \equiv \mathcal{J}^{(i)}_{kl} dA^{(k)} dA^{(i)} d\tilde{\tau}^{(i)} d\tau_{\delta} dp^{(k)}.$$

Furthermore, a computation shows that $\mathcal{J}^{(i)}_{kl} = |\eta|^{|c(i)|+d(k)+d(i)-1} \mathcal{J}^{(i)}_{kl} \circ \rho^{(i)}_{\delta}.$

**Second reduction.** Now, the group $G_{p^{(k)}}$ acts on $\gamma^{(k)}((S_{kl})_{p^{(k)}})$ with the isotropy types $(H_i) = (H_i_1), (H_i_{i+1}), \ldots, (H_i_L)$. By the same arguments given in the first reduction, the isotropy types occuring in $W_{kl}$ constitute a subset of these types, and we shall denote them by

$$(H_i) = (H_{i_1}), (H_{i_2}), \ldots, (H_{i_L}).$$

Consequently, for each $p^{(k)} \in M_k(H_k)$, $G_{p^{(k)}}$ acts on $S_{kl}$ with the isotropy types $(H_{i_1}), \ldots, (H_{i_L})$. If the isotropy groups $H_{i_1}, \ldots, H_{i_L}$ have the same dimensions, we shall see that the singularities of $\gamma^{(i)}((N)\setminus(N)^{-1})$ can be locally resolved over $Z_{kl}$. Moreover, if $G_{p^{(k)}}$ acts on $S_{kl}$ only with type $(H_L)$, the ideal $I_p$ can be partially monomialized in such a way that the critical sets of the corresponding weak transforms are clean. But since this is not the case in general, we have to continue with the iteration.

**N-th decomposition.** Denote by $\Lambda \leq L$ the maximal number of elements that a totally ordered subset of the set of isotropy types can have. Assume that $3 \leq N < \Lambda$, and let $\{(H_i_1), \ldots, (H_{i_N})\}$ be a totally ordered subset of the set of isotropy types such that $i_1 < \cdots < i_N < L$. Let $f_{i_1}, f_{i_1,i_2}, S_{i_1}, S_{i_1,i_2}$, as well as $p^{(i_2)} \in M_k(H_{i_1})$, $p^{(i_2)} \in \gamma^{(1)}((S_{i_1}^+)_{p^{(i_1)}})_{i_2}(H_{i_2}), \ldots$ be defined as in the first two iteration steps, and assume that $f_{i_1}, S_{i_1}, \ldots, p^{(i)}$, $\ldots$ have already been defined for $j < N$. For every fixed $p^{(i_{N-1})}$, let $\gamma^{(i_{N-1})}((S_{i_1}^{(i_1)})(\cdots_{i_{N-1}}))_{i_N}$ be the submanifold with corners of the closed $G_{p^{(i_{N-1})}}$-manifold $\gamma^{(i_{N-1})}((S_{i_1}^{(i_1)})(\cdots_{i_{N-1}}))_{i_N}$ from which all orbit types less than $G/H_{i_N}$ have been removed. Consider the invariant tubular neighborhood $f_{i_1}, i_N = \exp \circ \gamma^{(i_N)}: \nu_{i_1}, i_N \to \gamma^{(i_{N-1})}((S_{i_1,i_{N-1}})^{p^{(i_{N-1})}})_{i_N}$

of the set $\gamma^{(i_{N-1})}((S_{i_1,i_{N-1}})^{p^{(i_{N-1})}})_{i_N}$, and define $S_{i_1,i_{N-1}}$ as the sphere subbundle in $\nu_{i_1,i_N}$, while

$S_{i_1,i_{N-1}} = \left\{ v \in S_{i_1,i_{N-1}} : v = \sum v_i v_i^{(i_1,i_{N-1})}, v_{i_{N-1}} > 0 \right\}$

for some $\varrho_{i_{N-1}}$. Put $W_{i_1,i_{N-1}} = f_{i_1,i_{N-1}}(D_1(\nu_{i_1,i_N}))$, and denote the corresponding integral in the decomposition of $I_{i_1}^{(i_{N-1})}((S_{i_1,i_{N-1}})^{p^{(i_{N-1})}})_{i_N}$ by $I_{i_1,i_{N-1}}^{(i_{N-1})}(\mu)$. Here we can assume that, modulo terms of lower order, the $W_{i_1,i_{N-1}} \times G_{p^{(i_{N-1})}}$-support of the integrand in $I_{i_1,i_{N-1}}^{(i_{N-1})}(\mu)$ is contained in a compactum of a tubular neighborhood of the induced bundle $\pi_{i_1,i_{N-1}}\circ \gamma^{(i_{N-1})}(\nu_{i_1,i_{N-1}})^{(i_{N-1})}_{i_N}(H_{i_N})$, where $\pi_{i_1,i_{N-1}} : W_{i_1,i_{N-1}} \to \gamma^{(i_{N-1})}(\nu_{i_1,i_{N-1}})^{(i_{N-1})}_{i_N}(H_{i_N}), \gamma^{(i_{N-1})}(\nu_{i_1,i_{N-1}})^{(i_{N-1})}_{i_N}(H_{i_N})$ denotes the canonical projection. For a given point $p^{(i_N)} \in \gamma^{(i_{N-1})}(\nu_{i_1,i_{N-1}})^{(i_{N-1})}_{i_N}(H_{i_N})$, consider further the decomposition

$g_{p^{(i_{N-1})}} = g_{p^{(i_N)}} \oplus g_{p^{(i_{N-1})}},$

and set $d^{(i_N)} = \dim g_{p^{(i_{N-1})}}$; $e^{(i_N)} = \dim g_{p^{(i_N)}}$. This yields the decomposition

$$g = g_{p^{(i_1)}} \oplus g_{p^{(i_2)}} \oplus g_{p^{(i_3)}} \oplus \cdots = g_{p^{(i_1)}} \oplus g_{p^{(i_2)}} \oplus \cdots \oplus g_{p^{(i_{N-1})}}.$$
Denote by \( \{A_{i}^{(1)}(p_{1}^{(1)}, \ldots, p_{i}^{(1)})\} \) a basis of \( g_{p_{i}^{(1)}} \), and by \( \{B_{i}^{(1)}(p_{i}^{(1)}, \ldots, p_{i}^{(1)})\} \) a basis of \( g_{p_{i}^{(1)}} \). For \( A^{(i)} \in g_{p_{i}^{(1)}} \) and \( B^{(i)} \in g_{p_{i}^{(1)}} \) write further

\[
A^{(1)} = \sum_{i=1}^{d^{(1)}} \alpha^{(1)}_{i} A_{i}^{(1)}(p_{1}^{(1)}, \ldots, p_{1}^{(1)}), \quad B^{(1)} = \sum_{i=1}^{d^{(1)}} \beta^{(1)}_{i} B_{i}^{(1)}(p_{1}^{(1)}, \ldots, p_{1}^{(1)}),
\]

and let \( \{v_{i}^{(1)}(p_{i}^{(1)}, \ldots, p_{i}^{(1)})\} \) be an orthonormal frame in \( (p_{i}^{(1)}, \ldots, p_{i}^{(1)}) \).

**N-th monoidal transformation.** Let the monoidal transformations \( \zeta^{(1)}, \zeta^{(2)} \) be defined as in the first two iteration steps, and assume that the monoidal transformations \( \zeta^{(j)} \) have already been defined for \( j < N \). Put \( \tilde{M}^{(j)} = B_{Z^{(j)}}(\tilde{M}^{(j-1)}) \), \( \tilde{M}^{(0)} = M = M \times G \), and consider the monoidal transformation

\[
(35) \quad \zeta^{(N)} : B_{Z^{(N)}}(\tilde{M}^{(N-1)}) \rightarrow \tilde{M}^{(N-1)}, \quad Z^{(N)} = \bigcup_{i_{1} < \cdots < i_{N} < L} Z_{i_{1} \cdots i_{N}}, \quad (\text{disjoint union}),
\]

where the union is over all totally ordered subsets \( \{(H_{i_{1}}, \ldots, H_{i_{N}})\} \) of \( N \) elements with \( i_{1} < \cdots < i_{N} < L \), and

\[
Z_{i_{1} \cdots i_{N}} \simeq \bigcup_{p_{i_{1}}, \ldots, p_{i_{N}}} (-1, 1)^{N-1} \times \text{Iso}(i_{1} \cdots i_{N})^{-1}(S_{i_{1} \cdots i_{N}})^{(i_{1} \cdots i_{N})})_{i_{1}}(H_{i_{1}})
\]

are the possible maximal singular loci of \( (\zeta^{(1)} \circ \cdots \circ \zeta^{(N-1)})^{-1}(N) \). Denote by \( \zeta^{(1)}_{i_{1}} \circ \cdots \circ \zeta^{(N)}_{i_{N}} \) a local realization of the sequence of monoidal transformations \( \zeta^{(1)} \circ \cdots \circ \zeta^{(N)} \) corresponding to the totally ordered subset \( \{(H_{i_{1}}, \ldots, H_{i_{N}})\} \) in a set of \( (\theta^{(i_{1})}, \ldots, \theta^{(i_{N})}) \)-charts labeled by the indices \( i_{1}, \ldots, i_{N} \). As a consequence, we obtain a partial monomialization of the inverse image ideal sheaf \( (\zeta^{(1)} \circ \cdots \circ \zeta^{(N)})^{*}(I_{N}) \) according to

\[
(\zeta^{(1)} \circ \cdots \circ \zeta^{(N)})^{*}(I_{N}) \cdot E_{m, N}^{(N)} = \tau_{i_{1}} \cdots \tau_{i_{N}} \cdot (\zeta^{(1)} \circ \cdots \circ \zeta^{(N)})_{*}^{-1}(I_{N}) \cdot E_{m, N}^{(N)}
\]

in a neighborhood of any point \( \tilde{m} \in \tilde{M}^{(N)} = B_{Z^{(N)}}(\tilde{M}^{(N-1)}) \), as well as local factorizations of the phase function according to

\[
\Phi \circ (\zeta^{(1)} \circ \cdots \circ \zeta^{(N)})^{*}(I_{N}) \cdot \mathcal{E}_{m, N}^{(N)} = \tau_{i_{1}} \cdots \tau_{i_{N}} \cdot (\zeta^{(1)} \circ \cdots \circ \zeta^{(N)})^{*}(I_{N}) \cdot \mathcal{E}_{m, N}^{(N)}
\]

where in the relevant \( (\theta^{(i_{1})}, \ldots, \theta^{(i_{N})}) \)-charts

\[
\tilde{\Phi}^{(i_{1})} = \bigcup_{j=1}^{N} \partial(\tau_{i_{j}} A_{i}^{(1)}), \quad \tilde{\Phi}^{(i_{1})} = \bigcup_{j=1}^{N} \partial(\tau_{i_{j}} A_{i}^{(1)}), \quad \tilde{\Phi}^{(i_{1})} = \bigcup_{j=1}^{N} \partial(\tau_{i_{j}} A_{i}^{(1)}), \quad \tilde{\Phi}^{(i_{1})} = \bigcup_{j=1}^{N} \partial(\tau_{i_{j}} A_{i}^{(1)}), \quad \tilde{\Phi}^{(i_{1})} = \bigcup_{j=1}^{N} \partial(\tau_{i_{j}} A_{i}^{(1)}),
\]

the \( \{p_{i}^{(j)}\} \) being local coordinates, \( \tilde{\Phi}^{(i_{1})}(p_{i}^{(1)}, \theta^{(i_{1})}) \in \gamma^{(i_{1})}((S_{i_{1} \cdots i_{N}})^{(i_{1} \cdots i_{N})})_{i_{1}}(H^{(i_{1})}) \), and

\[
\Xi^{(i_{1})} = \Xi^{(i_{1})}, \quad \Xi^{(i_{1})} = \Xi^{(i_{1})}, \quad \Xi^{(i_{1})} = \Xi^{(i_{1})}, \quad \Xi^{(i_{1})} = \Xi^{(i_{1})}, \quad \Xi^{(i_{1})} = \Xi^{(i_{1})}, \quad \Xi^{(i_{1})} = \Xi^{(i_{1})},
\]

the Jacobian of the coordinate change given by

\[
\theta^{(i)} = \theta^{(i)}(\tau_{i_{j}} \exp_{p_{i_{j}+1}} \gamma_{i_{j}+1}(p_{i_{j}+1})).
\]
Modulo lower order terms, $I(\mu)$ is then given by a sum of integrals of the form

$$I_{i_1 \ldots i_N}^{\theta_1 \ldots \theta_N}(\mu) =$$

$$\int_{M_1(H_1) \times (-1,1)} \left[ \int_{\gamma(i_1)}((S_{i_1})_{p(i_1)})_2(H_2) \times (-1,1) \cdots \int_{\gamma(i_N)}((S_{i_1 \ldots i_N})_{p(i_N)})_N(H_N) \times (-1,1) \right] e^{i\mu \tau_1 \ldots \tau_N} \left( \frac{1}{\pi} \int_{\gamma(i_1)}\cdots \int_{\gamma(i_N)} \sum_{\tau = 1}^{i_N} d\tau \right) d\tau_1 \ldots d\tau_N \prod_{i = 1}^{N} \left[ \gamma(i) + \sum_{r = 1}^{j} d^{(r)} \right]$$

(36)

Here $a_{i_1 \ldots i_N}^{\theta_1 \ldots \theta_N}$ are amplitudes with compact support in a system of $(\theta^{(i_1)}, \ldots, \theta^{(i_N)})$-charts labeled by the indices $\theta_{i_1}, \ldots, \theta_{i_N}$, while

$$\gamma_{i_1 \ldots i_N}^{\theta_1 \ldots \theta_N} = \prod_{j=1}^{N} \left[ \gamma_j(e^{(i)}) + \sum_{r=1}^{j} d^{(r)} - 1 \right] \gamma_{i_1 \ldots i_N}^{\theta_1 \ldots \theta_N},$$

where $\gamma_{i_1 \ldots i_N}^{\theta_1 \ldots \theta_N}$ are functions which do not depend on the variables $\tau_{i_j}$.

**N-th reduction.** For each $p(i_{N-1})$, the isotropy group $G_{p(i_{N-1})}$ acts on $\gamma(i_{N-1})((S_{i_1 \ldots i_{N-1}})_{p(i_{N-1})})_{i_{N-1}}$ by the types $(H_{i_1}), \ldots, (H_{i_{N-1}})$. The types occurring in $W_{i_1 \ldots i_{N-1}}$ constitute a subset of these, and $G_{p(i_{N-1})}$ acts on the sphere bundle $S_{i_1 \ldots i_{N-1}}$ over the submanifold $\gamma(i_{N-1})((S_{i_1 \ldots i_{N-1}})_{p(i_{N-1})})_{i_{N-1}}(H_{i_{N-1}}) \subset W_{i_1 \ldots i_{N-1}}$ with one type less.

**End of iteration.** As before, let $\Lambda \leq L$ be the maximal number of elements of a totally ordered subset of the set of isotropy types. After $N = \Lambda - 1$ steps, the end of the iteration is reached. In particular, we will have achieved a desingularization of $\mathcal{N}$. For this, it is actually sufficient to consider only monoidal transformations (35) whose centers $Z^{(N)}$ are unions over totally ordered subsets $\{(H_{i_1}), \ldots, (H_{i_N})\}$ for which the corresponding orbit types $G/H_{i_j}$ are singular.

**Theorem 6.** Consider a compact, connected $n$-dimensional Riemannian manifold $M$, together with a compact, connected Lie groups $G$ acting effectively and isometrically on $M$, and put

$$\mathcal{N} = \{(x, g) \in M : gx = x\},$$

For every $1 \leq N \leq \Lambda - 1$, let the monoidal transformation $\zeta^{(N)}$ be defined as in (35), where $Z^{(N)}$ is an union over totally ordered subsets $\{(H_{i_1}), \ldots, (H_{i_N})\}$ of singular isotropy types of $N$ elements. Denote the sequence of monoidal transformations $\zeta^{(1)} \circ \cdots \circ \zeta^{(\Lambda - 1)}$ by $\zeta$, and put $\mathcal{M} = \mathcal{M}^{(\Lambda - 1)}$. Then $\zeta : \mathcal{M} \rightarrow \mathcal{M}$ yields a strong resolution of $\mathcal{N}$.

**Proof.** If all $G$-orbits on $M$ have the same dimension, $\mathcal{N}$ is a manifold, and $\zeta : \mathcal{M} \rightarrow \mathcal{M}$ is the identity. Let us therefore assume that there are singular orbits, and begin by recalling the covering

$$\mathcal{N} = \mathcal{N}_1 \cup \cdots \cup \mathcal{N}_L,$$

where $\mathcal{N}_L = \text{Iso } W_L$ is a manifold, and the $\mathcal{N}_k = \mathcal{N} \cap U_k$ are in general singular for $k < L$. Let $1 \leq N \leq \Lambda - 1$, and consider a totally ordered subset $\{(H_{i_1}), \ldots, (H_{i_N})\}$ of isotropy types such that $i_1 < \cdots < i_N$. In case that $(H_{i_1})$ is exceptional, all types $(H_{i_j})$ with $j < j'$ are exceptional, or principal. Indeed, if $H_{i_1}/H_L$ is finite and non-trivial, $H_{i_j}/H_L$ is also finite. In particular, if $(H_{i_1})$ is exceptional, $\mathcal{N}_{i_1}$ is a manifold. In what follows, let us therefore restrict to the case where $\{(H_{i_1}), \ldots, (H_{i_N})\}$ is a totally ordered subset of singular isotropy types which is maximal in the sense that there is no singular isotropy type $(H_{i_{N+1}})$ with $i_N < i_{N+1}$ such that $\{(H_{i_1}), \ldots, (H_{i_{N+1}})\}$ is a totally ordered subset. Let $\zeta^{(i_1)} \circ \cdots \circ \zeta^{(i_N)}$ be a local realization.
of the sequence of monoidal transformations $\zeta^{(1)} \circ \cdots \circ \zeta^{(N)}$ corresponding to the totally ordered subset $\{ (H_{i_1}), \ldots, (H_{i_N}) \}$ in a set of $(\theta^{(i_1)}, \ldots, \theta^{(i_N)})$-charts labeled by the indices $\theta_{i_1}, \ldots, \theta_{i_N}$. The preimage of $N_{i_1}$ under $\zeta^{(1)}_i \circ \cdots \circ \zeta^{(N)}_i$ is given by all points

$$(\tau_{i_1}, \ldots, \tau_{i_N}, p^{(i_1)}, \ldots, p^{(i_N)}, \tilde{\nu}^{(i_1)}, \ldots, A^{(i_1)}, \ldots, A^{(i_N)}, h^{(i_N)})$$

satisfying

$$(x^{(i_1 \cdots i_N)}, g^{(i_1 \cdots i_N)}) \in \mathcal{N},$$

where for $j = 1, \ldots, N$ we set

\begin{align*}
(x^{(i_1 \cdots i_N)}) &= \exp_{p^{(i_j)}}[\tau_{i_1} \exp_{p^{(i_{j+1})}}[\tau_{i_{j+1}} \exp_{p^{(i_{j+2})}}[\cdots \exp_{p^{(i_{N-1})}}[\tau_{i_{N-2}} \exp_{p^{(i_N)}}[\tau_{i_N} \zeta^{(i_N)}]]]]]]], \\
(g^{(i_1 \cdots i_N)} &= \exp_{\tau_{i_1} \cdots \tau_{i_{j-1}} A^{(i_j)}} \exp_{\tau_{i_j+1} \cdots \tau_{i_N} A^{(i_{j+1})}} \cdots \exp_{\tau_{i_{N-1}} \cdots \tau_{i_N} A^{(i_{N-1})}} \exp_{\tau_{i_N} A^{(i_N)}} h^{(i_N)}).
\end{align*}

Assume now that $\tau_{i_1} \cdots \tau_{i_N} \neq 0$. Since the point $x^{(i_1 \cdots i_N)}$ lies in a slice around $G \cdot p^{(i_j)}$, the condition $g^{(i_1 \cdots i_N)} \in G_x^{(i_1 \cdots i_N)}$ implies that $g^{(i_1 \cdots i_N)}$ must stabilize $p^{(i_1)}$ as well. From the inclusions

\begin{equation}
G_{p^{(i_N)}} \subset G_{p^{(i_{N-1})}} \subset \cdots \subset G_{p^{(i_1)}},
\end{equation}

and $g^{(i_1 \cdots i_N)}_{p^{(i_{j+1})}} \subset g^{(i_1 \cdots i_N)}_{p^{(i_j)}}$ one deduces $g^{(i_1 \cdots i_N)} \in G_{p^{(i_1)}}$, and we obtain

$$g^{(i_1 \cdots i_N)} p^{(i_j)} = e^{\tau_{i_1} \cdots \tau_{i_N} A^{(i_1)}} \cdot \sigma^{(i_j)} A^{(i_1)} p^{(i_1)} = p^{(i_1)}.$$ 

Thus we conclude $\sigma^{(i_j)} = 0$, which implies $g^{(i_1 \cdots i_N)} \in G_x^{(i_1 \cdots i_N)}$, and consequently $g^{(i_1 \cdots i_N)} \in G_x^{(i_1 \cdots i_N)}$. Repeating the above argument we see that

$$(x^{(i_1 \cdots i_N)}, g^{(i_1 \cdots i_N)}) \in \mathcal{N} \iff A^{(i_1)} = 0, \quad h^{(i_N)} \in \mathcal{E}^{(i_N)}$$

in case that $\tau_{i_1} \cdots \tau_{i_N} \neq 0$. Actually we have shown that if $\tau_{i_1} \cdots \tau_{i_N} \neq 0$

\begin{equation}
G_x^{(i_1 \cdots i_N)} = \mathcal{E}^{(i_N)},
\end{equation}

since $G_{\mathcal{E}^{(i_N)}} \subset G_{p^{(i_N)}}$. The preimage of $N_{i_1}$ under $\zeta^{(1)}_{i_1} \circ \cdots \circ \zeta^{(N)}_{i_1}$ is therefore given by

$$\left\{ \tau_{i_1} \cdots \tau_{i_N} \neq 0, \quad A^{(i_1)} = 0, \quad h^{(i_N)} \in \mathcal{E}^{(i_N)} \right\} \cup \bigcup_{j=1}^N \left\{ \tau_{i_j} = 0 \right\}.$$ 

By assumption, $G_{p^{(i_N)}}$ acts on $(S_{i_1 \cdots i_N})_{p^{(i_N)}}$ with orbits of the same dimension, so that

\begin{equation}
\{ A^{(i_1)} = 0, \quad h^{(i_N)} \in \mathcal{E}^{(i_N)} \}
\end{equation}

is a smooth submanifold, being equal to the total space of the isotropy bundle given by the local trivialization

$$(\tau_{i_j}, p^{(i_1)}, h^{(i_N)}, \mathcal{E}^{(i_N)}_{p^{(i_N)}}) \mapsto (\tau_{i_j}, p^{(i_1)}, \tilde{\nu}^{(i_N)}).$$

Now, for $1 \leq N \leq \Lambda - 1$, let $\zeta^{(N)}$ be defined as in [34], where $\mathcal{Z}^{(N)}$ is a union over totally ordered subsets of singular isotropy types of $N$ elements, and put $\zeta = \zeta^{(1)} \circ \cdots \circ \zeta^{(\Lambda-1)}$. By construction, $\zeta$ is given locally by sequences of local transformations $\zeta^{(1)}_{i_1} \circ \cdots \circ \zeta^{(N)}_{i_1}$ corresponding to maximal, totally ordered subsets $\{ (H_{i_1}), \ldots, (H_{i_N}) \}$ of singular isotropy types of $N \leq \Lambda - 1$ elements. Taking the union over all the corresponding sets [35], yields a smooth submanifold $\mathcal{N}$ which has normals crossings with the exceptional divisor $\zeta^{-1}(\text{Sing} \mathcal{N}) \subset \mathcal{M}$. Furthermore, $\zeta$ maps the union of the sets $\{ \tau_{i_1} \cdots \tau_{i_N} \neq 0, \quad A^{(i_1)} = 0, \quad h^{(i_N)} \in \mathcal{E}^{(i_N)} \}$ bijectively onto the non-singular part $\text{Reg} \mathcal{N}$ of $\mathcal{N}$. However, $\text{Reg} \mathcal{N}$ is not necessarily dense in $\mathcal{N}$, nor is $\zeta(\mathcal{N})$, so that $\zeta : \mathcal{N} \to \mathcal{N}$ might not be a birational map in general. Nevertheless, by successively blowing up the intersections of $\mathcal{N}$ with $\zeta^{-1}(\text{Sing} \mathcal{N})$ one finally obtains a strong resolution of $\mathcal{N}$.

\[\square\]
The resolution of $N$ constructed in Theorem 6 was deduced from a monomialization of the ideal sheaf $I_N$

$$
\zeta^*(I_N) \cdot \mathcal{E}_{\bar{m},A} = \tau_{i_1} \cdots \tau_{i_{k-1}} \cdot \zeta_{i_k}^{-1}(I_N) \cdot \mathcal{E}_{\bar{m},A}, \quad \bar{m} \in \bar{M},
$$

where $\zeta_{i_k}^{-1}(I_N)$ is a resolved ideal sheaf. In the following two sections, we shall derive from this a partial monomialization of the local ideal $I_\Phi = (\Phi)$ such that the corresponding weak transforms of $\Phi$ have clean critical sets. This will allow us to derive asymptotics for the integrals $I_{i_1 \cdots i_N}^B(\mu)$ in Section 8 via the stationary phase theorem.

6. Phase analysis of the weak transforms. The first main theorem

We continue with the notation of the previous sections and recall that the sequence of monoidal transformations $\zeta = \zeta^{(1)} \cdots \zeta^{(N-1)}$ is given locally by sequences of local transformations $\zeta_{i_1}^{(1)} \circ \cdots \circ \zeta_{i_N}^{(N)}$ corresponding to totally ordered subsets $\{(H_{i_1}), \ldots, (H_{i_N})\}$ of non-principal isotropy types that are maximal in the sense that, if there is an isotropy type $(H_{i_{N+1}})$ with $i_N < i_{N+1}$ such that $\{(H_{i_1}), \ldots, (H_{i_{N+1}})\}$ is a totally ordered subset, then $H_{i_{N+1}} = (H_L)$. Let now $x \in M$ be fixed, and $Z_x \subset T_x M$ be a neighborhood of zero such that $\exp_x : Z_x \rightarrow M$ is a diffeomorphism onto its image. One has

$$(\exp_x)_{*, v} : T_x Z_x \rightarrow T_{\exp_x v} M, \quad v \in Z_x,$$

and under the identification $T_x M \simeq T_0 Z_x$, one computes $(\exp_x)_{*, 0} \equiv \text{id}$. Furthermore, for $g \in G$ we have $g \cdot \exp_x v = L_g(\exp_x v) = \exp_{L_g(x)}(L_g)_{*, x}(v)$. Consider next a maximal, totally ordered subset $\{(H_{i_1}), \ldots, (H_{i_N})\}$ of isotropy types with $i_1 < \cdots < i_N < L$, and denote by

$$
\lambda : g_{p^{(1)}} \rightarrow g(\nu_{i_1, p^{(1)}}), \quad B^{(1)} \rightarrow \frac{d}{dt}(L_{e^{-itB^{(1)}}})_{*, p^{(1)}}|_{t=0},
$$

the linear representation of $g_{p^{(1)}}$ in $\nu_{i_1, p^{(1)}}$, where $\nu_{i_1, p^{(1)}} \in M_{i_1}(H_{i_1})$. For an arbitrary element $A^{(i_j)} \in g_{p^{(1)}}$ with $2 \leq j \leq N$, and $x^{(i_1 \cdots i_N)}$ given as in 37, one computes

$$
(A^{(i_j)})_{x^{(i_1 \cdots i_N)}} = \frac{d}{dt} \exp_{p^{(1)}} \left[ \tau_{i_1} \exp_{p^{(2)}} \left[ \cdots \left[ \tau_{i_{j-1}} \left( e^{-tA^{(i_j)}} \right)_{*, p^{(1)}} x^{(i_j \cdots i_N)} \right] \right] \right]|_{t=0},
$$

successively obtaining

$$
(A^{(i_j)})_{x^{(i_1 \cdots i_N)}} = \frac{d}{dt} \exp_{p^{(1)}} \left[ \tau_{i_1} \exp_{p^{(2)}} \left[ \cdots \left[ \tau_{i_{j-1}} \left( e^{-tA^{(i_j)}} \right)_{*, p^{(1)}} x^{(i_j \cdots i_N)} \right] \right] \right]|_{t=0},
$$

where we made the canonical identification $T_x (\nu_{i_1, p^{(1)}}) \equiv \nu_{i_1, p^{(1)}}$ for any $v \in (\nu_{i_1, p^{(1)})}$. We shall next define certain geometric distributions $E^{(i_j)}$ and $F^{(i_N)}$ on $M$ by setting

$$
E^{(i_1 \cdots i_N)} = \text{Span} \{ \tilde{Y}_{x^{(i_1 \cdots i_N)}} : Y \in g_{p^{(1)}} \},
$$

$$
E^{(i_1 \cdots i_N)} = \left( \exp_{p^{(1)}} \right)_{*, \tau_{i_1} x^{(i_2 \cdots i_N)}} \left( \exp_{p^{(2)}} \right)_{*, \tau_{i_2} x^{(i_3 \cdots i_N)}} \left( \exp_{p^{(3)}} \right)_{*, \tau_{i_3} x^{(i_4 \cdots i_N)}} \left[ \lambda(g_{p^{(1)}}) x^{(i_j \cdots i_N)} \right],
$$

where $2 \leq j \leq N$. By construction, if $\tau_{i_1} \cdots \tau_{i_N} \neq 0$, the $G$-orbit through $x^{(i_1 \cdots i_N)}$ is of principal type $G/H_L$, which amounts to the fact that $g_{p^{(N-1)}}$ acts on $S_{i_1 \cdots i_N}$ only with isotropy type $(H_L)$,
where we understand that $G_{p^{(i_0)}} = G$. Furthermore, (33) and (41) imply that

$$T_{x^{(i_1...i_N)}(G \cdot x^{(i_1...i_N)})} = E_{x^{(i_1...i_N)}}^{(i_1)} \oplus \bigoplus_{j=2}^{N} \tau_{i_j} \cdots \tau_{i_2} \tau_{i_1} E_{x^{(i_1...i_N)}}^{(i_2)} \oplus \tau_{i_1} \cdots \tau_{i_N} E_{x^{(i_1...i_N)}}^{(i_N)}.$$  

The main result of this section is the following

**Theorem 7** (First Main Theorem). Let $\{(H_{i_1}),..., (H_{i_N})\}$ be a maximal, totally ordered subset of non-principal isotropy types, and $\zeta^{(i_1)}, \cdots, \zeta^{(i_N)}$ a corresponding local realization of the sequence of monoidal transformations $\zeta^{(1)} \circ \cdots \circ \zeta^{(N)}$ in a set of $\{\theta^{(i_1)}, \cdots, \theta^{(i_N)}\}$-charts labeled by the indices $\sigma_{i_1}, \cdots, \sigma_{i_N}$. Consider the corresponding factorization

$$\Phi \circ (id_1 \otimes \cdots \otimes \zeta^{(i_1)} \circ \cdots \circ \zeta^{(i_N)}) = (\iota_{1...i_N}) \tilde{\Phi}_{\text{tot}} = \tau_{i_1} \cdots \tau_{i_N} (\iota_{1...i_N}) \tilde{\Phi}_{\text{wk, pre}}$$

of the phase function (41) where

$$(\iota_{1...i_N}) \tilde{\Phi}_{\text{wk, pre}} = \left( \Xi, T (dp^{(i_1)}_1 (A^{(i_1)}_{p^{(i_1)}_1}), \cdots, dp^{(i_2)}_1 (A^{(i_2)}_{p^{(i_2)}_1}), \cdots, 0, dp^{(i_N)}_1 (A^{(i_N)}_{p^{(i_N)}_1}), \cdots), \theta^{(i_1)}_1 \left( \tilde{\omega}^{(i_1)}_1 \right) - \theta^{(i_1)}_1 \left( (h^{(i_1)}_{p^{(i_1)}_1})_{\ast, p^{(i_1)}_1} \tilde{\omega}^{(i_1)}_1 \right), \cdots, \right) + \sum_{j=1}^{N} O(\|\tau_{i_j} A^{(i_1)}_j\|).$$

Let further $(\iota_{1...i_N}) \tilde{\Phi}_{\text{wk}}$ denote the pullback of $(\iota_{1...i_N}) \tilde{\Phi}_{\text{wk, pre}}$ along the substitution $\tau = \delta_{i_1...i_N}(\sigma)$ given by the sequence of local quadratic transformations

$$\delta_{i_1...i_N}: (\sigma_{i_1}, \cdots, \sigma_{i_N}) \mapsto (\sigma_{i_1}', \sigma_{i_2}, \cdots, \sigma_{i_N}) = (\sigma_{i_1}', \cdots, \sigma_{i_1}'', \sigma_{i_2}', \cdots, \sigma_{i_N}).$$

Then the critical set Crit$(\iota_{1...i_N}) \tilde{\Phi}_{\text{wk}}$ of $(\iota_{1...i_N}) \tilde{\Phi}_{\text{wk}}$ is given by all points

$$(\sigma_{i_1}, \cdots, \sigma_{i_N}, p^{(i_1)}, \cdots, p^{(i_N)}, \tilde{\omega}^{(i_1)}, A^{(i_1)}, \cdots, A^{(i_N)}, h^{(i_N)})(\xi)$$

satisfying the conditions

(I) $A^{(i_j)} = 0$ for all $j = 1, \cdots, N$, and $(h^{(i_N)}_{p^{(i_N)}_1})_{\ast, p^{(i_N)}_1} \tilde{\omega}^{(i_N)} = \tilde{\omega}^{(i_N)}$,

(II) $\eta_{x^{(i_1...i_N)}} \in \text{Ann}(E_{x^{(i_1...i_N)}}^{(i_j)})$ for all $j = 1, \cdots, N,$

(III) $\eta_{x^{(i_1...i_N)}} \in \text{Ann}(E_{x^{(i_1...i_N)}}^{(i_1)}),$

where $\eta$ denotes the 1-form $\sum_{i=1}^{n} \xi_i d\tilde{x}_i$. Furthermore, Crit$(\iota_{1...i_N}) \tilde{\Phi}_{\text{wk}}$ is a $C^{\infty}$-submanifold of codimension $2\kappa$, where $\kappa = \dim \overline{G/H_L}$ is the dimension of a principal orbit.

**Proof.** In what follows, set

$$Z^{\delta_{i_1...i_N}} = (\zeta^{\delta_{i_1...i_N}} \circ \cdots \circ \zeta^{\delta_{i_1...i_N}} \circ (\delta_{i_1...i_N} \otimes id_\xi)) \otimes id_\xi$$

so that

$$\Phi \circ Z^{\delta_{i_1...i_N}} = (\iota_{1...i_N}) \tilde{\Phi}_{\text{tot}} = \tau_{i_1} \cdots \tau_{i_N}(\sigma) (\iota_{1...i_N}) \tilde{\Phi}_{\text{wk}},$$

and let $\sigma_{i_1} \cdots \sigma_{i_N} \neq 0$. In this case, $Z^{\delta_{i_1...i_N}}$ constitutes a diffeomorphism, so that

$$\text{Crit}(\iota_{1...i_N}) \tilde{\Phi}_{\text{tot}} |_{\sigma_{i_1} \cdots \sigma_{i_N} \neq 0} = \{ (\sigma_{i_1}, \cdots, \sigma_{i_N}, p^{(i_1)}, \cdots, p^{(i_N)}, \tilde{\omega}^{(i_1)}, A^{(i_1)}, \cdots, A^{(i_N)}, h^{(i_N)})(\xi) : (x^{(i_1...i_N)})(\xi, g^{(i_1...i_N)}) \in C, \sigma_{i_1} \cdots \sigma_{i_N} \neq 0 \}.$$ 

\[\text{Note that } (\iota_{1...i_N}) \tilde{\Phi}_{\text{wk, pre}} \text{ was formerly denoted by } (\iota_{1...i_N}) \tilde{\Phi}_{\text{wk}}.\]
where we employed the notation of \[37\]. Now, by \[21\],

\[
(x^{(i_1\ldots i_N)}, \xi, g^{(i_1\ldots i_N)}) \in \mathcal{C} \iff \eta_{x^{(i_1\ldots i_N)}} = \sum_{i=1}^n \xi_i (d\hat{x}_i)_{x^{(i_1\ldots i_N)}} \in \Omega,
\]

\[
g^{(i_1\ldots i_N)} \in G_{x^{(i_1\ldots i_N)}}, \eta_{x^{(i_1\ldots i_N)}}).
\]

The reasoning which led to \[39\] in particular implies that condition (I) is equivalent to \(g^{(i_1\ldots i_N)} \in G_{x^{(i_1\ldots i_N)}}\) in case that all \(\sigma_{ij}\) are different from zero. Now, \(\eta_{x^{(i_1\ldots i_N)}} \in \Omega\) means that

\[
\sum \xi_i (d\hat{x}_i)_{x^{(i_1\ldots i_N)}} \in \text{Ann}(T_{\hat{x}^{(i_1\ldots i_N)}}(G \cdot x^{(i_1\ldots i_N)})).
\]

But if \(\sigma_{ij} \neq 0\) for all \(j = 1, \ldots, N\), (II) and (III) imply that

\[
\eta_{x^{(i_1\ldots i_N)}}((\exp_{p^{(i_1)}})_{\tau_1} x^{(i_1\ldots i_N)}[\ldots (\exp_{p^{(i_j)}})_{\tau_j} x^{(i_1\ldots i_N)}[\lambda(Z)x^{(i_N)}]\ldots]) = 0 \quad \forall Z \in g_{p^{(i_{N-1})}},
\]

since \(g_{p^{(i_{N-1})}} = g_{p^{(i_{N})}} \oplus g_{p^{(i_{N})}}^{-1}\). By repeatedly using this argument, we conclude that under the assumption \(\sigma_{i_1} \cdots \sigma_{i_N} \neq 0\)

\[
(45) \quad (\text{II), (III)} \iff \eta_{x^{(i_1\ldots i_N)}} \in \text{Ann}(T_{\hat{x}^{(i_1\ldots i_N)}}(G \cdot x^{(i_1\ldots i_N)})).
\]

Taking everything together therefore gives

\[
\text{Crit}((i_1\ldots i_N)\hat{\Phi}^{tot})_{\sigma_{i_1}\cdots\sigma_{i_N} \neq 0} = \{(\sigma_{i_1}, \ldots, \sigma_{i_N}, p^{(i_1)}, \ldots, p^{(i_N)}, \tilde{v}^{(i_N)}, A^{(i_1)}, \ldots, A^{(i_N)}, h^{(i_N)}, \xi) : \sigma_{i_1} \cdots \sigma_{i_N} \neq 0, (\text{I)-(III)} \text{ are fulfilled and } h^{(i_N)} \cdot \eta_{x^{(i_1\ldots i_N)}} = \eta_{x^{(i_1\ldots i_N)}}\}.
\]

and we assert that

\[
\text{Crit}((i_1\ldots i_N)\hat{\Phi}^{wk}) = \text{Crit}((i_1\ldots i_N)\hat{\Phi}^{tot})_{\sigma_{i_1}\cdots\sigma_{i_N} \neq 0}.
\]

To show this, let us still assume that all \(\sigma_{ij}\) are different from zero. Then all \(\tau_{ij}\) are different from zero, too, and \(\partial_{\tilde{\Phi}} (i_1\ldots i_N)\hat{\Phi}^{wk} = 0\) is equivalent to

\[
\partial_{\tilde{\Phi}} \Phi(x^{(i_1\ldots i_N)}, \xi, g^{(i_1\ldots i_N)}) = 0,
\]

which gives us the condition \(g^{(i_1\ldots i_N)} \in G_{x^{(i_1\ldots i_N)}}\). By the reasoning which led to \[39\] we therefore obtain condition (I) in the case that all \(\sigma_{ij}\) are different from zero. Let now one of the \(\sigma_{ij}\) be equal to zero. Then all \(\tau_{ij}\) are zero, too, and \(\partial_{\tilde{\Phi}} (i_1\ldots i_N)\hat{\Phi}^{wk} = 0\) is equivalent to

\[
(47) \quad \tilde{A}^{(i_j)}_{p^{(i_j)}}(p^{(i_j)}) = 0 \quad \text{for all } 1 \leq j \leq N \text{ and } q, \quad (1 - h^{(i_N)})(p^{(i_N)}) \tilde{v}^{(i_N)} = 0
\]

since the \((n \times n)\)-matrix \(\Xi\) is invertible, so that the kernel of the corresponding linear transformation is trivial. Denote by \(N_{p^{(i_1)}}(G \cdot p^{(i_1)})\) the normal space in \(T_{p^{(i_1)}}M\) to the orbit \(G \cdot p^{(i_1)}\), on which \(G_{p^{(i_1)}}\) acts, and define \(N_{p^{(i_1)}}(G \cdot p^{(i_1)} \cdot p^{(i_1+1)})\) successively as the normal space to the orbit \(G_{p^{(i_1)}} \cdot p^{(i_1+1)}\) in the \(G_{p^{(i_1)}}\)-space \(N_{p^{(i_1)}}(G \cdot p^{(i_1)} \cdot p^{(i_1)})\), where we understand that \(G_{p^{(0)}} = G\). Since smooth actions of compact Lie groups are locally smooth, the aforementioned actions can be assumed to be orthogonal, see [8], pages 171 and 308. Since \(\tilde{A}^{(i_j)}_{p^{(i_j)}} \in T_{p^{(i_j)}}(G \cdot p^{(i_1)})\) is tangent to \(M_j(H_{ij})\), and \(\tilde{A}^{(i_j)}_{p^{(i_j)}} \in T_{p^{(i_j)}}(G \cdot p^{(i_1)} \cdot p^{(i_1)})\) is tangent to \(\gamma^{(i_1)}((S_{i_1\ldots i_{j-1}}^{+})_p^{(i_1-1)}) \cdot H_{ij}\), we finally obtain

\[
(48) \quad \partial_{\tilde{\Phi}} (i_1\ldots i_N)\hat{\Phi}^{wk} = 0 \iff (\text{I})
\]

for arbitrary \(\sigma_{ij}\). In particular, one concludes that \((i_1\ldots i_N)\hat{\Phi}^{wk}\) must vanish on its critical set. Since

\[
d(', '(i_1\ldots i_N)\hat{\Phi}^{tot}) = d('', 'i_1\ldots i_1)\hat{\Phi}^{wk} + \tau_{i_1} \cdots \tau_{i_N} d('', 'i_1\ldots i_N)\hat{\Phi}^{wk},
\]
one sees that
\[ \text{Crit}(\sigma_1 \cdots \sigma_N \Phi_{\text{tot}}) \subset \text{Crit}(\sigma_1 \cdots \sigma_N \Phi_{\text{tot}}). \]

In turn, the vanishing of \( \Phi \) on its critical set implies
\[ \text{Crit}(\sigma_1 \cdots \sigma_N \Phi_{\text{tot}}) \subset \text{Crit}(\sigma_1 \cdots \sigma_N \Phi_{\text{tot}}). \]

Therefore, by continuity,
\[ \text{Crit}(\sigma_1 \cdots \sigma_N \Phi_{\text{tot}}) \subset \text{Crit}(\sigma_1 \cdots \sigma_N \Phi_{\text{tot}}). \]

In order to see the converse inclusion we shall henceforth assume that \( \partial_{x^k} \Phi_{\text{tot}} = 0 \), and consider next the \( \alpha \)-derivatives, where we shall again take \( \sigma_i = \sigma_i \neq 0 \). Taking into account (39) and (48), one sees that
\[ \partial_{x^k} \Phi_{\text{tot}} = 0 \quad \iff \quad \sum_{q=1}^N \xi_q (A_r^{(i)})(x_{\sigma_1 \cdots \sigma_N}) = 0. \]

By (41) we therefore obtain for arbitrary \( \sigma \) and \( 1 \leq j \leq N \)
\[ \partial_{x^k} \Phi_{\text{tot}} = 0 \quad \iff \quad \sum_{q=1}^N \xi_q (A_r^{(i)})(x_{\sigma_1 \cdots \sigma_N}) = 0. \]

Consequently,
\[ \partial_{x^k} \Phi_{\text{tot}} = 0 \quad \iff \quad \alpha_i = \alpha_i \neq 0 \quad (\text{II}). \]

In a similar way, one sees that
\[ \partial_{x^k} \Phi_{\text{tot}} = 0 \quad \iff \quad \beta_i = \beta_i \neq 0 \quad (\text{III}). \]

by which the necessity of the conditions (I)–(III) is established. In order to see their sufficiency, let them be fulfilled, and assume again that \( \sigma_i = \sigma_i \neq 0 \) for all \( j = 1, \ldots, N \). Then (16) implies that
\[ \eta_{x^{(i_1 \cdots i_N)}} \in \text{Ann}(T_{x^{(i_1 \cdots i_N)}}(G \cdot x^{(i_1 \cdots i_N)})). \]

Now, if \( \sigma_i \cdots \sigma_i \neq 0 \), \( G \cdot x^{(i_1 \cdots i_N)} \) is of principal type \( G/H_L \) in \( M \), so that the isotropy group of \( x^{(i_1 \cdots i_N)} \) must act trivially on \( N_{x^{(i_1 \cdots i_N)}}(G \cdot x^{(i_1 \cdots i_N)}) \), compare [8], page 181. If therefore \( \mathcal{X} = \mathcal{X}_T + \mathcal{X}_N \) denotes an arbitrary element in \( T_{x^{(i_1 \cdots i_N)}}(G \cdot x^{(i_1 \cdots i_N)}) \), and \( g \in G_{x^{(i_1 \cdots i_N)}} \), one computes
\[ \eta_{x^{(i_1 \cdots i_N)}}(\mathcal{X}_N) = \eta_{x^{(i_1 \cdots i_N)}}(\mathcal{X}_N). \]

With (39) we then conclude that \( h^{(i_1 \cdots i_N)} \cdot \eta_{x^{(i_1 \cdots i_N)}} = \eta_{x^{(i_1 \cdots i_N)}} \), since \( (h^{(i_1 \cdots i_N)})_{x^{(i_1 \cdots i_N)}} \). Set next
\[ V^{(i_1 \cdots i_N)} = N_{p^{(i)}}(G_{p^{(i_j-1)}} \cdot p^{(i_j)}). \]

We then have the following

**Lemma 2.** The orbit of the point \( \tilde{v}^{(i_1 \cdots i_N)} \) in the \( G_{p^{(i_1 \cdots i_N)}} \)-space \( V^{(i_1 \cdots i_N)} \) is of principal type.

**Proof of the lemma.** By assumption, for \( \sigma_i = \sigma_i \neq 0 \), \( 1 \leq j \leq N \), the \( G \)-orbit of \( x^{(i_1 \cdots i_N)} \) is of principal type \( G/H_L \) in \( M \). The theory of compact group actions then implies that this is equivalent to the fact that \( x^{(i_1 \cdots i_N)} \in V^{(i_1 \cdots i_N)} \) is of principal type in the \( G_{p^{(i_1 \cdots i_N)}} \)-space \( V^{(i_1 \cdots i_N)} \), see [8], page 181, which in turn is equivalent to the fact that \( x^{(i_1 \cdots i_N)} \in V^{(i_1 \cdots i_N)} \) is of principal type in the \( G_{p^{(i_1 \cdots i_N)}} \)-space \( V^{(i_1 \cdots i_N)} \), and so forth. Thus, \( x^{(i_1 \cdots i_N)} \in V^{(i_1 \cdots i_N)} \) must be of principal type in the \( G_{p^{(i_1 \cdots i_N)}} \)-space \( V^{(i_1 \cdots i_N)} \) for all \( j = 1, \ldots, N \), and the assertion follows. \( \square \)
Let us now assume that one of the $\sigma_{i_j}$ vanishes. Then

\begin{equation}
(\text{II), (III)} \iff \begin{cases} 
\eta_{p(i_j)} \in \text{Ann}(E_{p(i_j)}^{(i_j)}) & \forall j = 1, \ldots, N, \\
\eta_{p(i_j)} \in \text{Ann}(F_{p(i_j)}^{(i_N)}) \end{cases},
\end{equation}

where $E_{p(i_j)}^{(i_1)} = T_{p(i_1)}(G \cdot p(i_1))$, and

\begin{equation}
E_{p(i_j)}^{(i_j)} \simeq T_{p(i_j)}(G_{p(i_j-1)} \cdot p(i_j)) \subset V^{(i_1 \ldots i_{j-1})}, \quad 2 \leq j \leq N,
\end{equation}

while $F_{p(i_j)}^{(i_N)} \simeq T_{p(i_j)}(G_{p(i_j N)} \cdot \tilde{v}^{(i_N)}) \subset V^{(i_1 \ldots i_N)}$. Consequently, we obtain the direct sum of vector spaces

$$E_{p(i_1)}^{(i_1)} \oplus E_{p(i_1)}^{(i_2)} \oplus \cdots \oplus E_{p(i_1)}^{(i_N)} \oplus F_{p(i_1)}^{(i_N)} \subset T_{p(i_1)}M.$$ 

Now, as a consequence of the previous lemma, the stabilizer of $\tilde{v}^{(i_N)}$ must act trivially on $N_{\tilde{v}^{(i_N)}}(G_{p(i_N)} \cdot \tilde{v}^{(i_N)})$. If therefore $T = T_N + T_N$ denotes an arbitrary element in

$$T_{p(i_1)}M \simeq \bigoplus_{j=1}^{N} T_{p(i_j)}(G_{p(i_j-1)} \cdot p(i_j)) \oplus T_{\tilde{v}^{(i_N)}}(G_{p(i_N)} \cdot \tilde{v}^{(i_N)}) \oplus N_{\tilde{v}^{(i_N)}}(G_{p(i_N)} \cdot \tilde{v}^{(i_N)})$$

and $G_{\tilde{v}^{(i_N)}} \subset G_{p(i_N)}$ imply that for $g \in G_{\tilde{v}^{(i_N)}}$

$$g \cdot \eta_{p(i_1)}(\tilde{x}) = [(L_{g^{-1}})^*_{p(i_1)} \eta_{p(i_1)}(\tilde{x})] = \eta_{p(i_1)}(L_{g^{-1}})_*(\tilde{x}_N) \sim \eta_{p(i_1)}(\tilde{x}_N) = \eta_{p(i_1)}(\tilde{x}).$$

Collecting everything together we have shown for arbitrary $\sigma = (\sigma_{i_1}, \ldots, \sigma_{i_N})$ that

\begin{equation}
\partial_{\xi_{i_1}, \ldots, \xi_{i_N}}(i_1 \ldots i_N) \eta_{x^{(i_1 \ldots i_N)}} \tilde{\Phi}^{\text{wk}} = 0 \iff \text{ (I), (II), (III)} \implies h^{(i_N)} \in \mathcal{G}_{x^{(i_1 \ldots i_N)}}.
\end{equation}

By (10) and (11) we therefore conclude

\begin{equation}
\text{Crit}(\tilde{\Phi}^{\text{tot}}|_{\sigma_{i_1} \ldots \sigma_{i_N} \neq 0} = \text{Crit}(\tilde{\Phi}^{\text{wk}}).
\end{equation}

We have thus computed the critical set of $(\tilde{\Phi}^{\text{wk}}$, and it remains to show that it is a $C^\infty$- submanifold of codimension $2\kappa$. By the previous considerations,

$$\text{Crit}(\tilde{\Phi}^{\text{wk}})$$

\begin{equation}
= \left\{ A^{(i_1)} = 0, \quad h^{(i_N)} \in \mathcal{G}_{\tilde{v}^{(i_N)}}, \quad \eta_{x^{(i_1 \ldots i_N)}} \in \text{Ann}(\bigoplus_{j=1}^{N} E_{x^{(i_1 \ldots i_N)}}^{(i_1)} \oplus F_{x^{(i_1 \ldots i_N)}}^{(i_N)}) \right\}.
\end{equation}

Now, since for $\sigma_{i_1} \cdots \sigma_{i_N} \neq 0$ the $G$-orbit of $x^{(i_1 \ldots i_N)}$ is of principal type $G/H_L$ in $M$, (13) implies in this case that

$$\kappa = \dim T_{x^{(i_1 \ldots i_N)}}(G \cdot x^{(i_1 \ldots i_N)}) = \dim \left[ E_{x^{(i_1 \ldots i_N)}}^{(i_1)} \oplus \bigoplus_{j=2}^{N} \tau_{i_1} \cdots \tau_{i_{j-1}} E_{x^{(i_1 \ldots i_N)}}^{(i_j)} \oplus \tau_{i_1} \cdots \tau_{i_N} F_{x^{(i_1 \ldots i_N)}}^{(i_N)} \right]$$

$$= \sum_{j=1}^{N} \dim \left[ E_{x^{(i_1 \ldots i_N)}}^{(i_j)} \oplus \dim F_{x^{(i_1 \ldots i_N)}}^{(i_N)} \right].$$
But \( \dim E_{x}^{(i_{1}, \ldots, i_{N})} = \dim G^{i_{j-1}} \cdot p^{(i_{j})} \) in particular shows that the dimensions of the spaces \( E_{x}^{(i_{1}, \ldots, i_{N})} \) do not depend on the variables \( \sigma_{i_{j}} \). A similar argument applies to \( F_{x}^{(i_{1}, \ldots, i_{N})} \), so that we obtain the equality

\[
\kappa = \sum_{j=1}^{N} \dim E_{x}^{(i_{1}, \ldots, i_{N})} + \dim F_{x}^{(i_{1}, \ldots, i_{N})}
\]

for arbitrary \( x^{(i_{1}, \ldots, i_{N})} \). Note that, in contrast, the dimension of \( T_{x}^{(i_{1}, \ldots, i_{N})}(G \cdot x^{(i_{1}, \ldots, i_{N})}) \) collapses, as soon as one of the \( \tau_{i_{j}} \) becomes zero. Since the annihilator of a subspace of \( T_{x}M \) is a linear subspace of \( T_{x}^{\ast}M \), we arrive at a vector bundle with \( E \) given by the trivialization

\[
\bigg( \sigma_{i_{j}}, p^{(i_{j})}, \bar{v}^{(i_{N})}, \text{Ann} \left( \bigoplus_{j=1}^{N} E_{x}^{(i_{1}, \ldots, i_{N})} \oplus F_{x}^{(i_{1}, \ldots, i_{N})} \right) \bigg) \rightarrow \big( \sigma_{i_{j}}, p^{(i_{j})}, \bar{v}^{(i_{N})} \big).
\]

Consequently, by (53) and (58) we see that \( \text{Crit}(\bar{v}^{(i_{1}, \ldots, i_{N})}) \) is equal to the total space of the fiber product of the mentioned vector bundle with the isotropy bundle given by the local trivialization

\[
\big( \sigma_{i_{j}}, p^{(i_{j})}, \bar{v}^{(i_{N})}, G_{\bar{v}^{(i_{N})}} \big) \rightarrow \big( \sigma_{i_{j}}, p^{(i_{j})}, \bar{v}^{(i_{N})} \big).
\]

Lastly, equation (39) implies \( \dim g_{\bar{v}^{(i_{N})}} = d - \kappa \), which concludes the proof of Theorem 7. 

7. Phase analysis of the weak transforms. The second main theorem

In this section, we shall prove that the Hessians of the weak transforms \( \bar{v}^{(i_{1}, \ldots, i_{N})} \) are transversally non-degenerate at each point of their critical sets. We begin with the following general observation. Let \( M \) be a \( n \)-dimensional \( C^{\infty} \)-manifold, and \( C \) the critical set of a function \( \psi \in C^{\infty}(M) \), which is assumed to be a smooth submanifold in a chart \( \mathcal{O} \subset M \). Let further

\[
\alpha : (x, y) \mapsto m, \quad \beta : (q_{1}, \ldots, q_{n}) \mapsto m, \quad m \in \mathcal{O},
\]

be two systems of local coordinates on \( \mathcal{O} \), such that \( \alpha(x, y) \in C \) if and only if \( y = 0 \). One computes

\[
\partial_{y_{i}}(\psi \circ \alpha)(x, y) = \sum_{i=1}^{n} \frac{\partial(\psi \circ \beta)}{\partial q_{i}}(\beta^{-1} \circ \alpha(x, y)) \partial_{y_{i}}(\beta^{-1} \circ \alpha)_{i}(x, y),
\]

as well as

\[
\partial_{y_{k}} \partial_{y_{l}}(\psi \circ \alpha)(x, y) = \sum_{i=1}^{n} \frac{\partial(\psi \circ \beta)}{\partial q_{i}}(\beta^{-1} \circ \alpha(x, y)) \partial_{y_{k}} \partial_{y_{l}}(\beta^{-1} \circ \alpha)_{i}(x, y) + \sum_{i,j=1}^{n} \frac{\partial^{2}(\psi \circ \beta)}{\partial q_{i} \partial q_{j}}(\beta^{-1} \circ \alpha(x, y)) \partial_{y_{k}}(\beta^{-1} \circ \alpha)_{j}(x, y) \partial_{y_{l}}(\beta^{-1} \circ \alpha)_{i}(x, y).
\]

Since

\[
\alpha_{*}(x, y)(\partial_{y_{k}}) = \sum_{j=1}^{n} \partial_{y_{k}}(\beta^{-1} \circ \alpha)_{j}(x, y) \beta_{*}(\beta^{-1} \circ \alpha)_{j}(x, y)(\partial_{q_{j}}),
\]

this implies

\[
\partial_{y_{k}} \partial_{y_{l}}(\psi \circ \alpha)(x, 0) = \text{Hess} \psi_{\alpha(x, 0)}(\alpha_{*}(x, 0)(\partial_{y_{k}}), \alpha_{*}(x, 0)(\partial_{y_{l}})),
\]

by definition of the Hessian (33), Section 2. Let us now write \( x = (x', x'') \), and consider the restriction of \( \psi \) onto the \( C^{\infty} \)-submanifold \( M_{c'} = \{ m \in \mathcal{O} : m = \alpha(c', x'', y) \} \). We write \( \psi_{c'} = \psi_{|M_{c'}} \),
and denote the critical set of $\psi_{c'}$ by $C_{c'}$, which contains $C \cap M_{c'}$ as a subset. Introducing on $M_{c'}$ the local coordinates $\alpha' : (x'', y) \mapsto (c', x'', y)$, we obtain

$$\partial_{y_k} \partial_{y_l} (\psi_{c'} \circ \alpha')(x'', 0) = \text{Hess} \psi_{c'} \circ (\partial_{x''} \alpha')(0) \partial_{y_k} \partial_{y_l} \alpha'(x'', 0).$$

Let us now assume $C_{c'} = C \cap M_{c'}$, a transversal intersection. Then $C_{c'}$ is a submanifold of $M_{c'}$, and the complement of $T_{\alpha'(x'', 0)} C_{c'}$ in $T_{\alpha'(x'', 0)} M_{c'}$ at a point $\alpha'(x'', 0)$ is spanned by the vector fields $\alpha'_{*,(x'', 0)}(\partial_{y_k})$. Since clearly

$$\partial_{y_k} \partial_{y_l} (\psi_{c'} \circ \alpha')(x'', 0) = \partial_{y_k} \partial_{y_l} (\psi \circ \alpha)(x, 0), \quad x = (c', x''),$$

we thus have proven the following

**Lemma 3.** Assume that $C_{c'} = C \cap M_{c'}$. Then $\psi$ is transversally non-degenerate at $\alpha'(c', x'', 0) \in C$ if, and only if $\text{Hess} \psi_{c'}$ is transversally non-degenerate at $\alpha'(x'', 0) \in C_{c'}$. That is, $\text{Hess} \psi$ defines a non-degenerate quadratic form on $T_{\alpha'(x'', 0)} M / T_{\alpha'(x'', 0)} C_{c'}$ if, and only if $\text{Hess} \psi_{c'}$ defines a non-degenerate quadratic form on $T_{\alpha'(x'', 0)} M_{c'} / T_{\alpha'(x'', 0)} C_{c'}$.

Let us now state the main result of this section, the notation being the same as in the previous sections.

**Theorem 8** (Second Main Theorem). Let $\{(H_{i_1}, \ldots, (H_{i_N})\}$ be a maximal, totally ordered subset of non-principal isotropy types of the $G$-action on $M$, and $Z^{\tilde{t}_{i_1} \cdots \tilde{t}_{i_N}}_{\tilde{t}_{i_1} \cdots \tilde{t}_{i_N}}$ be defined as in (44). Consider the corresponding factorization

$$\Phi \circ Z^{\tilde{t}_{i_1} \cdots \tilde{t}_{i_N}}_{\tilde{t}_{i_1} \cdots \tilde{t}_{i_N}} \tilde{Phi}_{\text{tot}} = \tau_{i_1}(\sigma) \cdots \tau_{i_N}(\sigma) \tilde{Phi}_{\text{wk}}$$

of the phase function (10). Then, at each point of the critical manifold $\text{Crit}(\tilde{Phi}_{\text{tot}})$, the Hessian $\text{Hess} \tilde{Phi}_{\text{tot}}$ is transversally non-degenerate.

For the proof of Theorem 8 we need the following

**Lemma 4.** Let $(x, \xi, g) \in \text{Crit}(\Phi)$, and $x \in M(H_L)$. Then $(x, \xi, g) \in \text{Reg Crit}(\Phi)$, and $\text{Hess} \Phi$ is transversally non-degenerate at $(x, \xi, g)$.

**Proof.** The first assertion is clear from (15) - (17). To see the second, consider the 1-form $\eta = \sum \xi_i d\tilde{x}_i$, and note that by (16)

$$\eta_x \in \Omega \cap T_x Y, \quad x \in M(H_L), \quad g \in G_x \implies g \cdot \eta_x = \eta_x.$$ 

Since by (18) the condition $d_x \Phi(x, \xi, g) = 0$ is equivalent to $g \cdot \eta_x = \eta_x$, and

$$d_x \Phi(x, \xi, g) = 0 \iff g_x = x, \quad d_g \Phi(x, \xi, g) = 0 \iff \eta_{gx} \in \Omega,$$

we obtain on $T^*(Y \cap M(H_L)) \times G$ the implication

$$d_{\xi, g} \Phi(x, \xi, g) = 0 \implies d_x \Phi(x, \xi, g) = 0.$$

Let $\Phi_2(\xi, g)$ denote the phase function (10) regarded as a function of the coordinates $\xi, g$ alone, while $x$ is regarded as a parameter. Lemma 3 then implies that on $T^*(Y \cap M(H_L)) \times G$ the study of the transversal Hessian of $\Phi$ can be reduced to the study of the transversal Hessian of $\Phi_2$. Now, with respect to the coordinates $\xi, g$, the Hessian of $\Phi_2$ is given by

$$\begin{pmatrix}
0 & (d\tilde{x}_i)_x((\tilde{X}_i)_x) \\
(d\tilde{x}_i)_x((\tilde{X}_i)_x) & (\partial_{X_i, X_j} \Phi_2)(\xi, g)
\end{pmatrix}.$$
where \(\{X_1, \ldots, X_d\}\) denotes a basis of \(\mathfrak{g}\). A computation then shows that the kernel of the corresponding linear transformation is equal to \( \left\{ (\xi, \tilde{s}) \in \text{Ann} (T_x (G \cdot x)) : \sum \tilde{s}_j (\tilde{X}_j)_x = 0 \right\} \simeq T_{\xi, g} (\text{Crit} (\Phi_x)). \) The lemma now follows by the following general observation. Let \(B\) be a symmetric bilinear form on an \(n\)-dimensional \(\mathbb{K}\)-vector space \(V\), and \(B = (B_{ij})_{i,j}\) the corresponding Grammian matrix with respect to a basis \(\{v_1, \ldots, v_n\}\) of \(V\) such that

\[
B(u, w) = \sum_{i,j} u_i w_j B_{ij}, \quad u = \sum u_i v_i, \quad w = \sum w_i v_i.
\]

We denote the linear operator given by \(B\) with the same letter, and write

\[
V = \ker B \oplus W.
\]

Consider the restriction \(B_{|W \times W}\) of \(B\) to \(W \times W\), and assume that \(B_{|W \times W} (u, w) = 0\) for all \(u \in W\), but \(w \neq 0\). Since the Euclidean scalar product in \(V\) is non-degenerate, we necessarily must have \(Bw = 0\), and consequently \(w \in \ker B \cap W = \{0\}\), which is a contradiction. Therefore \(B_{|W \times W}\) defines a non-degenerate symmetric bilinear form.

**Proof of Theorem** \([\text{Theorem} \, 8]\) For \(\sigma_1, \ldots, \sigma_{i_N} \neq 0\), the sequence of monoidal transformations \(Z_{\sigma_1 \ldots \sigma_{i_N}}^{\Phi_{ij}}\) defined in \([\text{[4]}]\) is a diffeomorphism, so that by the previous lemma

\[
\text{Hess}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}} (\sigma_{ij}, p(i), \tilde{\alpha}(i), \tilde{\alpha}(i), h(i), \xi) = \text{transversally non-degenerate at each point of } \text{Crit}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}} \neq 0.
\]

Next, one computes

\[
\left( \frac{\partial^2 (i_1 \ldots i_N) \Phi_{\sigma_{ij}}}{\partial \gamma_k \partial \gamma_l} \right)_{k,l} = \tau_1 (\sigma) \cdots \tau_{i_N} (\sigma) \left( \frac{\partial^2 (i_1 \ldots i_N) \Phi_{\sigma_{ij}}}{\partial \gamma_k \partial \gamma_l} \right)_{k,l} + \left( \begin{array}{c}
\frac{\partial^2 (\tau_1 (\sigma) \cdots \tau_{i_N} (\sigma))}{\partial \sigma_r \sigma_s} \\
0
\end{array} \right)_{r,s}
\]

where \(\gamma_k\) stands for any of the coordinates, and \(R\) represents a matrix whose entries contain first order derivatives of \((i_1 \ldots i_N) \Phi_{\sigma_{ij}}\) as factors. But since by \([\text{[9]}]\)

\[
\text{Crit}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}} \neq 0 = \text{Crit}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}}|_{\sigma_{ij} \cdots \sigma_{i_N} \neq 0},
\]

we conclude that the transversal Hessian of \((i_1 \ldots i_N) \Phi_{\sigma_{ij}}\) does not degenerate along the manifold \(\text{Crit}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}}|_{\sigma_{ij} \cdots \sigma_{i_N} \neq 0}\). Therefore, it remains to study the transversal Hessian of \((i_1 \ldots i_N) \Phi_{\sigma_{ij}}\) in the case that any of the \(\sigma_{ij}\) vanishes, that is, along the exceptional divisor. Now, the proof of the first main theorem, in particular \([\text{[5]}]\), showed that

\[
\partial_{\xi, \alpha}(i_1, \ldots, \alpha(i_N), h(i_N)) (i_1 \ldots i_N) \Phi_{\sigma_{ij}} = 0 \implies \partial_{\sigma_{i_1}, \ldots, \sigma_{i_N}, p(i_1), \ldots, p(i_N), \tilde{\alpha}(i_N)} (i_1 \ldots i_N) \Phi_{\sigma_{ij}} = 0.
\]

If therefore

\[
(i_1 \ldots i_N) \Phi_{\sigma_{ij}}^{\sigma_{ij}, p(i), h(i)} (\alpha(i), h(i), \xi)
\]

denotes the weak transform of the phase function \(\Phi\) regarded as a function of the variables \((\alpha(i_1), \ldots, \alpha(i_N), h(i_N), \xi)\) alone, while the variables \((\sigma_{i_1}, \ldots, \sigma_{i_N}, p(i_1), \ldots, p(i_N), \tilde{\alpha}(i_N))\) are kept fixed at constant values,

\[
\text{Crit}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}, p(i), \tilde{\alpha}(i)} = \text{Crit}^{(i_1 \ldots i_N)} \Phi_{\sigma_{ij}} \cap \left\{ \sigma_{ij} , p(i), \tilde{\alpha}(i) = \text{constant} \right\}.
\]
Thus, the critical set of \((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\) is equal to the fiber over \((\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)})\) of the vector bundle
\[
\left((\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}), G_{\tilde{\varphi}^{(s)}} \times \text{Ann} \left( \bigoplus_{j=1}^N E_s^{(s_j)} \oplus E_s^{(s_1 \cdots s_j)} \right) \right) \mapsto (\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}),
\]
and in particular a smooth submanifold. Lemma 3 then implies that the study of the transversal Hessian of \((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\) can be reduced to the study of the transversal Hessian of \((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\).

The crucial fact is now contained in the following

**Proposition 3.** Assume that \(\sigma_{s_j} \cdots \sigma_{s_N} = 0\). Then
\[
\ker \text{Hess} \left((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\right)(0, \ldots, 0, h^{(s_1)}(x), \xi) \simeq T_{(0, \ldots, 0, h^{(s_1)}(x), \xi)} \text{Crit} \left((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\right)
\]
for all \((0, \ldots, 0, h^{(s_1)}(x), \xi) \in \text{Crit} \left((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\right),\) and arbitrary \(p^{(s_j)}, \tilde{\varphi}^{(s_j)}\).

**Proof.** Since \(\sigma_{s_j} \cdots \sigma_{s_N} = 0\), we have
\[
(s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}} = \left\{(t^{(s_1)}(A_{s_1}^{(s_1)}), \ldots, 0, dp^{(s_2)}(A_{s_2}^{(s_2)}), \ldots, 0, dp^{(s_N)}(\tilde{\varphi}^{(s_N)}), \ldots) = \right\}.
\]
The only non-vanishing second order derivatives at a critical point \((0, \ldots, 0, h^{(s_1)}(x), \xi)\) therefore read
\[
\begin{align*}
\frac{\partial}{\partial \sigma_{s_j}(s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}} &= \left\langle \Xi, T(0, \ldots, 0, dp^{(s_1)}(A_{s_1}^{(s_1)}), \ldots, 0, \xi) \right\rangle, \\
\frac{\partial}{\partial \beta^{(s_j)}(s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}} &= \left\langle \Xi, T(0, \ldots, 0, \theta^{(s_1)}(\lambda(B_{s_1}^{(s_1)})\lambda(B_{s_2}^{(s_2)})\tilde{\varphi}^{(s_2)}), \ldots, \xi) \right\rangle,
\end{align*}
\]
where we used canonical coordinates of the first kind on \(G_{p^{(s_1)}}\) of the form \(e^{\sum \alpha^{(s_1)}(p^{(s_1)}) \cdot h^{(s_1)}}\).

Consequently, the Hessian of the function \((s_1, \ldots, s_N)_{\sigma_{s_j}, p^{(s_j)}, \tilde{\varphi}^{(s_j)}}\) with respect to the coordinates \(\xi, \alpha^{(s_j)}, \beta^{(s_j)}\) is given on its critical set by the matrix
\[
(61) \begin{pmatrix}
\Xi & 0 \\
0 & 1_d
\end{pmatrix}
\begin{pmatrix}
0 & T \xi & E & F
\end{pmatrix}
\begin{pmatrix}
T \Xi & 0 \\
0 & 1_d
\end{pmatrix},
\]
where the \((n \times d)\)-matrix \(E\) is defined by
\[
E = \begin{pmatrix}
0 & 0 & 0 & \ldots & 0 \\
0 & 0 & 0 & \ldots & 0 \\
0 & 0 & dp^{(s_2)}(A_{s_2}^{(s_2)}) & \ldots & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
0 & 0 & 0 & \ldots & d\theta^{(s_1)}(B_{s_1}^{(s_1)})
\end{pmatrix},
\]
and the \((d \times d)\)-matrix \(F\) by
\[
F = \begin{pmatrix}
0 & 0 & \Xi \\
0 & -T(0, \ldots, 0, \theta^{(s_1)}(\lambda(B_{s_1}^{(s_1)})\lambda(B_{s_2}^{(s_2)})\tilde{\varphi}^{(s_2)}), \ldots, \xi)
\end{pmatrix}.
\]
Since \( \Xi \) is invertible, the kernel of the linear transformation corresponding to (61) is isomorphic to the kernel of the linear transformation defined by

\[
\begin{pmatrix}
0 \\
T \xi \\
\mathcal{E}
\end{pmatrix}
\]

which we shall now compute. Clearly, the column vector \( T(\tilde{\xi}, \tilde{\alpha}^{(i_1)}, \ldots, \tilde{\alpha}^{(i_N)}, \tilde{\beta}^{(i_N)}) \) lies in the kernel if and only if

(a) \[ \sum_s \tilde{\alpha}^{(i_1)}(\tilde{A}^{(i_1)}_{s})_{p^{(i_1)}}(\tilde{d}^{(i_1)}_{p^{(i_1)}}) = 0 \] for all \( 1 \leq j \leq N \), \( \sum_m \tilde{\beta}^{(i_N)}_{m}(\tilde{B}^{(i_N)}_{m})_{\tilde{e}^{(i_N)}}(\tilde{e}^{(i_N)}) = 0 \);

(b) \[ \sum_{r=1}^{n-\epsilon^{(i_1)}} \xi_r dp^{(i_1)}_{p^{(i_1)}}(T_{p^{(i_1)}}(G \cdot p^{(i_1)})) = 0; \]

(b) \[ \sum_{r=1}^{n-\epsilon^{(i_1)}} -c^{(i_1)} -1 \tilde{\xi}_{n-c^{(i_1)}+r+1} dp^{(i_1)}_{p^{(i_1)}}(T_{p^{(i_1)}}(G_{p^{(i_1)}} \cdot p^{(i_1)})) = 0 \] for all \( 2 \leq j \leq N \);

(c) \[ \sum_{r=1}^{n} \tilde{\xi}_{n-\epsilon^{(i_N)}+r} dp^{(i_N)}_{p^{(i_N)}}(T_{p^{(i_N)}}(G_{p^{(i_N)}} \cdot \tilde{v}^{(i_N)})) = 0. \]

In this case, the vector \( T(\tilde{\xi}^{\prime}, \tilde{\alpha}^{(i_1)}, \ldots, \tilde{\alpha}^{(i_N)}, \tilde{\beta}^{(i_N)}) \) lies in the kernel of (61), where \( \tilde{\xi}^{\prime} = \tilde{\xi} \cdot \Xi^{-1} \).

Let now \( E^{(i)} \), \( F^{(i)} \), and \( V^{(i_1 \ldots i_N)} \) be defined as in (62) and (63), and remember that we have the isomorphisms (55). For condition (a) to hold, it is necessary and sufficient that

\[ \tilde{\alpha}^{(i_1)} = 0, \quad 1 \leq j \leq N, \quad \sum_m \tilde{\beta}^{(i_N)}_{m}(\tilde{B}^{(i_N)}_{m})_{\tilde{e}^{(i_N)}}(\tilde{e}^{(i_N)}) = 0. \]

On the other hand, condition (b1) means that on \( T_{p^{(i_1)}}(G \cdot p^{(i_1)}) \) we have

\[ 0 = \sum_{r=1}^{n-\epsilon^{(i_1)}} \tilde{\xi}_r dp^{(i_1)}_{p^{(i_1)}} = \sum_{r=1}^{n-\epsilon^{(i_1)}} (\tilde{\xi}^{\prime} : \Xi) dp^{(i_1)}_{p^{(i_1)}} = \sum_{i=1}^{n} \tilde{\xi}_i d\tilde{x}_i, \]

where \( (\tilde{x}_1, \ldots, \tilde{x}_n) = (p_1^{(i_1)}, \ldots, p^{(i_1)}_{n-\epsilon^{(i_1)}}, \theta^{(i_1)}_1, \ldots, \theta^{(i_1)}_{\epsilon^{(i_1)}}) \) are the coordinates introduced in Section 5. Indeed,

\[ d\theta^{(i_1)}_s(A^{(i_1)}_{p^{(i_1)}}) = A^{(i_1)}_{p^{(i_1)}}(\theta^{(i_1)}_s) = 0 \quad \forall s = 1, \ldots, \epsilon^{(i_1)}, \quad A^{(i_1)} \in \mathfrak{g}_p^{+} \]

\( M_i \) (\( H_i \)) being \( G \)-invariant. Therefore \( \sum_{r=1}^{n} \tilde{\xi}^{\prime}(d\tilde{x}_i)_{p^{(i_1)}} \in \text{Ann}(E^{(i_1)}_{p^{(i_1)}}) \). In the same way, condition (b2) implies that on \( T_{p^{(i_1)}}(G_{p^{(i_1)}} \cdot p^{(i_1)}), \)

\[
0 = \sum_{r=1}^{c^{(i_1)}-1} (\tilde{\xi}^{\prime} : \Xi)_{n-c^{(i_1)}+1} dp^{(i_1)} = \sum_{s=1}^{n-c^{(i_1)}} (\tilde{\xi}^{\prime} : \Xi)_{n-c^{(i_1)}+s} dp^{(i_1)} + \sum_{r=1}^{c^{(i_1)}-1} (\tilde{\xi}^{\prime} : \Xi)_{n-c^{(i_1)}+1} dp^{(i_1)} + \sum_{r=1}^{c^{(i_1)}-1} (\tilde{\xi}^{\prime} : \Xi)_{n-c^{(i_1)}+1} dp^{(i_1)}
\]

where we put \( \tilde{\xi}^{\prime} = \tilde{\xi}^{\prime} : \Xi(\{i_1\}, \ldots, \Xi(\{i_1\}, \ldots, \{i_1\}, \ldots, \{i_1\}, \ldots, \{i_1\}, \ldots, \{i_1\}, \ldots, \{i_1\}) \). Hereby we expressed the differential forms \( d\theta^{(i_1)}_s \) in terms of the differential forms \( d\tau_{j-1}, dp^{(i_1)}_{p^{(i_1)}}, \) and \( d\theta^{(i_1)}_s \), and took into account that the forms \( d\tau_{j-1} \) and \( dp^{(i_1)}_{p^{(i_1)}} \) vanish on \( T_{p^{(i_1)}}(G_{p^{(i_1)}} \cdot p^{(i_1)}) \subset (v_{i_1 \ldots i_1})_{p^{(i_1)}} \). Now, if \( v \in (v_{i_1 \ldots i_1})_{p^{(i_1)}}, \) and \( \sigma(v) = \exp_{p^{(i_1)}}(v), \)

\[ v(p^{(i_1)}_{p^{(i_1)}}) = \frac{d}{dt} p^{(i_1)}_{p^{(i_1)}}(\sigma(v))_{t=0} = 0, \]
so that the forms $dp_{n}^{(i_j-1)}$ must vanish on $(\nu_{i_1\ldots i_{j-1}})^{p_{i_j-1}}$. In case that $3 \leq j \leq N$, the same argument shows that the forms $d\tau_{i_{j-2}}$ also vanish on $(\nu_{i_1\ldots i_{j-1}})^{p_{i_j-1}}$, and proceeding inductively this way, we see that conditions $(b_1)$ and $(b_2)$ are equivalent to
\[ \sum_{i=1}^{n} \hat{\xi}_i(d\hat{x}_i)_{p^{(i_1)}} \in \text{Ann}(E_{p^{(i_1)}}) \quad \forall \quad j = 1, \ldots, N. \]

Similarly, one deduces that condition $(c)$ is equivalent to
\[ \sum_{i=1}^{n} \hat{\xi}_i(d\hat{x}_i)_{p^{(i_1)}} \in \text{Ann}(F_{p^{(i_1)}}). \]

On the other hand, by (58),
\[ \frac{2}{\pi} \int_{(0,\ldots,0,h^{(i_1)},\xi)} \]
\[ T_{(0,\ldots,0,h^{(i_1)},\xi)} \text{Crit}(\{ (i_1,\ldots,i_N) \}, \{ \hat{a}^{w_k}_{\sigma_{ij}, \rho_{ij}^{(i_N)}} \}, \{ \tilde{\xi}^{(i_N)} \}) = \left\{ (\tilde{\alpha}^{(i_1)}, \ldots, \tilde{\alpha}^{(i_N)}, \tilde{\beta}^{(i_N)}, \tilde{\gamma}^{(i_1)}): \tilde{\alpha}^{(i_j)} = 0, \right. \]
\[ \sum_{m} \tilde{\beta}^{(i_{N})}_{m}(B_{m}^{(i_{N})}) \in g_{E^{(i_{N})}}, \sum_{i} \hat{\xi}_i(d\hat{x}_i)_{p^{(i_1)}} \in \text{Ann}(\bigoplus_{j=1}^{N} F_{p^{(i_1)}}), \]
and the proposition follows.

The previous proposition now implies that for $\sigma_1 \cdots \sigma_{i_N} = 0$, the Hessian of $\{ (i_1,\ldots,i_N) \}, \{ \hat{a}^{w_k}_{\sigma_{ij}, \rho_{ij}^{(i_N)}} \}, \{ \tilde{\xi}^{(i_N)} \}$ is transversally non-degenerate at each point $(0,\ldots,0,h^{(i_N)},\xi)$ of its critical set, and Theorem 8 follows with Lemma 3.

8. Asymptotics for the integrals $I_{i_1\ldots i_N}^{\theta_{i_1}\cdots \theta_{i_N}}(\mu)$

We are now in position to describe the asymptotic behavior of the integrals $I_{i_1\ldots i_N}^{\theta_{i_1}\cdots \theta_{i_N}}(\mu)$ defined in (36) using the stationary phase theorem. Let $\{(H_{i_1}), \ldots, (H_{i_N})\}$ be a maximal, totally ordered subset of non-principal isotropy group, and $\zeta^{\theta_{i_1} \cdots \theta_{i_N}}(H_{i_1}) \cdots \zeta^{\theta_{i_1} \cdots \theta_{i_N}}(H_{i_N})$ a corresponding local realization of the sequence of monoidal transformations $\xi^{(1)}(H_{i_1}) \cdots \xi^{(N)}(H_{i_N})$ in a set of $\{ (\theta^{(i_1)}, \ldots, \theta^{(i_N)}) \}$-charts labeled by the indices $\theta_{i_1}, \ldots, \theta_{i_N}$. Since the considered integrals are absolutely convergent, we can interchange the order of integration by Fubini, and write
\[
I_{i_1\ldots i_N}^{\theta_{i_1}\cdots \theta_{i_N}}(\mu) = \int_{(-1,1)^N} \sum_{j=1}^{N} |\tau_{i_j}|^{c^{(i_j)} + \sum_{r=1}^{j} d^{(i_r)} - 1} d\tau_{i_1} \cdots d\tau_{i_N},
\]
where we set
\[
I_{i_1\ldots i_N}^{\theta_{i_1}\cdots \theta_{i_N}}(\nu) = \int_{M_{i_1}(H_{i_1})} \sum_{j=1}^{N} |\tau_{i_j}|^{c^{(i_j)} + \sum_{r=1}^{j} d^{(i_r)} - 1} d\tau_{i_1} \cdots d\tau_{i_N},
\]
and
\[
\int_{j^{(i_N)}}^{(i_1)}(\{ (S_{i_1})^{(i_1)} \}_{p^{(i_1)}})_{i_2}(H_{i_2}) \cdots \int_{j^{(i_N-1)}}^{(i_1)}(\{ (S_{i_1\ldots i_{N-1}})_{p^{(i_{N-1})}} \}_{i_N}(H_{i_N})}
\]
\[ e^{i\nu^{(1)}(1)\cdots \nu^{(N)}(N)} \delta_{\nu^{(1)}(1)\cdots \nu^{(N)}(N)} d\xi dA^{(i_1)} \cdots dA^{(i_N)} dh^{(i_N)} dp^{(i_N)} \cdots dp^{(i_1)},
\]
according to the notation introduced in Theorem 7 and introduced the new parameter
\[
\nu = \mu \cdot \tau_{i_1} \cdots \tau_{i_N}.\]
Now, for a sufficiently small $\varepsilon > 0$ to be chosen later we define
\[
1 \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu) = \int_{\Pi_{j=1}^N (-1,1) \cap (-\varepsilon, \varepsilon)} \frac{J^\vartheta (\mu \cdot \tau_{I_1} \cdots \tau_{I_N})}{\prod_{j=1}^N |\tau_{I_j}|^{|c(I_j)| + \sum_{r=1}^j d(\tau_r) - 1} d\tau_{I_N} \ldots d\tau_{I_1}} d\tau_{I_1},
\]
\[
2 \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu) = \int_{(-\varepsilon, \varepsilon)^N} \frac{J^\vartheta (\mu \cdot \tau_{I_1} \cdots \tau_{I_N})}{\prod_{j=1}^N |\tau_{I_j}|^{|c(I_j)| + \sum_{r=1}^j d(\tau_r) - 1} d\tau_{I_N} \ldots d\tau_{I_1}} d\tau_{I_1}.
\]

**Lemma 5.** One has $c(I_j) + \sum_{r=1}^j d(\tau_r) - 1 \geq \kappa$ for arbitrary $j = 1, \ldots, N$, where $\kappa$ denotes the dimension of $G/H_L$.

**Proof.** We first note that $c(I_j) = \dim (\nu_{I_1 \cdots I_j})_{p(I_j)} \geq \dim G_{p(I_j)} \cdot x^{(I_j+1 \cdots I_N)} + 1$. Indeed, $(\nu_{I_1 \cdots I_j})_{p(I_j)}$ is an orthogonal $G_{p(I_j)}$-space, so that the dimension of the $G_{p(I_j)}$-orbit of $x^{(I_j+1 \cdots I_N)} \in \gamma(I_j)((S_{I_1 \cdots I_j}^+)^{-1} p(I_j))$ can be at most $c(I_j) - 1$. Now, under the assumption $\sigma_{I_1} \cdots \sigma_{I_N} \neq 0$, (44) and (41) imply
\[
T_{x^{(I_1 \cdots I_N)}} (x^{(I_j+1 \cdots I_N)}) \simeq T_{x^{(I_1 \cdots I_N)}} (x^{(I_1 \cdots I_N)})
\]
where the distributions $E(I_j), F(I_j)$ where defined in (42). On then computes
\[
\dim G_{p(I_j)} \cdot x^{(I_j+1 \cdots I_N)} = \sum_{l=j+1}^N \dim E(I_j)_{x^{(I_1 \cdots I_N)}} + \dim F(I_j)_{x^{(I_1 \cdots I_N)}},
\]
which implies
\[
c(I_j) \geq \sum_{l=j+1}^N \dim E(I_j)_{x^{(I_1 \cdots I_N)}} + \dim F(I_j)_{x^{(I_1 \cdots I_N)}} + 1.
\]

But since $E(I_j)_{x^{(I_1 \cdots I_N)}} \simeq G_{p(I_j)} \cdot p(I_j)$ for any $\sigma$, the last inequality holds for arbitrary $\sigma$, too. On the other hand, one has
\[
d(I_j) = \dim g_{p(I_j)}^\perp = \dim [\lambda(\vartheta_{p(I_j)}) \cdot p(I_j)] = \dim [\lambda(\vartheta_{p(I_j)}) \cdot x^{(I_j \cdots I_N)}] = \dim E(I_j)_{x^{(I_1 \cdots I_N)}}.
\]

The assertion now follows with (69). \qed

As a consequence of the lemma, we obtain for $2 \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu)$ the estimate
\[
\left| 2 \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu) \right| \leq C \int_{(-\varepsilon, \varepsilon)^N} \prod_{j=1}^N |\tau_{I_j}|^{|c(I_j)| + \sum_{r=1}^j d(\tau_r) - 1} d\tau_{I_N} \ldots d\tau_{I_1}
\]
\[
\leq C \int_{(-\varepsilon, \varepsilon)^N} \prod_{j=1}^N |\tau_{I_j}|^\kappa d\tau_{I_N} \ldots d\tau_{I_1} = \frac{2C}{\kappa + 1} \varepsilon^{N(\kappa + 1)}
\]

for some $C > 0$. Let us now turn to the integral $1 \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu)$. After performing the change of variables $\delta_{I_1 \cdots I_N}$ one obtains for $1 \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu)$ the expression
\[
\int_{\varepsilon < |\tau_{I_j}(\sigma)| < 1} \hat{I}_{I_1 \ldots I_N}^\vartheta (\mu \cdot \tau_{I_1}(\sigma) \cdots \tau_{I_N}(\sigma)) \prod_{j=1}^N |\tau_{I_j}(\sigma)|^{|c(I_j)| + \sum_{r=1}^j d(\tau_r) - 1} |\det D\delta_{I_1 \cdots I_N}(\sigma)| d\sigma,
\]
where $J_{i_1 \ldots i_N}^{\sigma_1 \ldots \sigma_N} (\nu)$ is defined by the right hand side of (63), with $(i_1 \ldots i_N) \tilde{\Phi}^{wk, pre}$ being replaced by 
$(i_1 \ldots i_N) \tilde{\Phi}^{wk, \sigma}$, which denotes the weak transform $(i_1 \ldots i_N) \tilde{\Phi}^{wk}$ as a function of the variables $p^{(i_j)}, \tilde{g}^{(i_N)},$ 
$\alpha^{(i_j)}, h^{(i_N)}, \xi$ alone, while the variables $\sigma = (\sigma_{i_1}, \ldots, \sigma_{i_N})$ are regarded as parameters.

**Theorem 9.** Let $\sigma = (\sigma_{i_1}, \ldots, \sigma_{i_N})$ be a fixed set of parameters. Then, for every $\tilde{N} \in \mathbb{N}$ there exists a constant $C_{\tilde{N}, (i_1 \ldots i_N) \tilde{\Phi}^{wk}} > 0$ such that 

$$
|J_{i_1 \ldots i_N}^{\sigma_1 \ldots \sigma_N} (\nu) - (2\pi |\nu|^{-1})^\tilde{N} \sum_{j=0}^{\tilde{N}-1} |\nu|^{-j} Q_j ((i_1 \ldots i_N) \tilde{\Phi}^{wk}_{\sigma}, a_{i_1 \ldots i_N} J_{i_1 \ldots i_N})| \leq C_{\tilde{N}, (i_1 \ldots i_N) \tilde{\Phi}^{wk}} |\nu|^{-\tilde{N}},
$$

with estimates for the coefficients $Q_j$, and an explicit expression for $Q_0$. Moreover, the constants $C_{\tilde{N}, (i_1 \ldots i_N) \tilde{\Phi}^{wk}}$ and the coefficients $Q_j$ have uniform bounds in $\sigma$.

**Proof.** In what follows, we regard $\tilde{M} (N)$ as a Riemannian manifold with the product metric induced by the Riemannian metrics on $\gamma^{(i_j-1)} ((S_{i_1 \ldots i_{j-1}} p^{(i_{j-1})})_{i_j}) (H_{i_j}), \gamma^{(i_N)} (\{S_{i_1 \ldots i_N} p^{(i_N)}\}, g^{(i_N)}_{\mu(i_N)}, G_{p^{(i_N)}},$ and $(-1, 1)^N$, and corresponding volume density. Similarly, $\tilde{X}$ carries a Riemannian structure, being a paracompact manifold. As a consequence of the main theorems, and Lemma 3 together with the observations preceding Proposition 3 the phase function $(i_1 \ldots i_N) \tilde{\Phi}^{wk}$ has a clean critical set for any value of the parameters $\sigma$. That is, in the relevant charts we have

- the critical set of $(i_1 \ldots i_N) \tilde{\Phi}^{wk}$ is a $C^{\infty}$-submanifold of codimension $2\kappa$ for arbitrary $\sigma$;

- the Hessian of $(i_1 \ldots i_N) \tilde{\Phi}^{wk}$ is transversally non-degenerate at each point of its critical set.

Thus, the necessary conditions for applying the principle of the stationary phase to the integral $J_{i_1 \ldots i_N}^{\sigma_1 \ldots \sigma_N} (\nu)$ are fulfilled, and we obtain the desired asymptotic expansion by Theorem 3. Note that the amplitude $a_{i_1 \ldots i_N}^{\sigma_1 \ldots \sigma_N}$ might depend on $\mu$, compare the expression for $O(\mu^{n-2})$ in Theorem 4 the dependence being of the form $a_{i_1}(t, \kappa_{i_1}(x), \mu_{i_1})$, where $a_{i_1} \in S^0_{phg}$ is a classical symbol of order 0. But since for large $|\eta|$  \[3\] this dependence does not interfere with the asymptotics. To see the existence of the uniform bounds, note that by Remark 4 we have

$$
C_{\tilde{N}, (i_1 \ldots i_N) \tilde{\Phi}^{wk}} \leq C'_{\tilde{N}} \sup_{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} \left\| \left( \text{Hess}^{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} - 1 \right) \right\|.
$$

But since by Lemma 3 the transversal Hessian 

$$
\text{Hess}^{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} |_{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} \sup_{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} \left\| \left( \text{Hess}^{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} - 1 \right) \right\|.
$$

is given by

$$
\text{Hess}^{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} |_{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} \sup_{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} \left\| \left( \text{Hess}^{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} - 1 \right) \right\|.
$$

we finally obtain the estimate

$$
C_{\tilde{N}, (i_1 \ldots i_N) \tilde{\Phi}^{wk}} \leq C'_{\tilde{N}} \sup_{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} \left\| \left( \text{Hess}^{(i_1 \ldots i_N) \tilde{\Phi}^{wk}} - 1 \right) \right\| \leq C_{\tilde{N}, (i_1 \ldots i_N) \tilde{\Phi}^{wk}}
$$

by a constant independent of $\sigma$. Similarly, one can show the existence of bounds of the form

$$
|Q_j ((i_1 \ldots i_N) \tilde{\Phi}^{wk}, a_{i_1 \ldots i_N} \tilde{\Phi}_{i_1 \ldots i_N})| \leq C_{j, (i_1 \ldots i_N) \tilde{\Phi}^{wk}},
$$

\[3\] Hereby we are taking into account that due to the presence of the factor $\Delta_{\nu} (\zeta_{i_1}(t, \kappa_{i_1}(x), \eta))$, $|\eta|$ is bounded away from zero.
with constants $\tilde{C}_{j_{i_{1},\ldots,i_{N}}}$ independent of $\sigma$.

**Remark 2.** Before going on, let us remark that for the computation of the integrals $I_{i_{1},\ldots,i_{N}}^\varphi(\nu) (\mu)$ it is only necessary to have an asymptotic expansion for the integrals $J_{i_{1},\ldots,i_{N}}^\varphi(\nu)$ in the case that $\sigma_{i_{2}} \cdots \sigma_{i_{N}} \neq 0$, which can also be obtained without the main theorems using only the factorization of the phase function $\Phi$ given by the resolution process, together with Lemma 4. Nevertheless, the main consequence to be drawn from the main theorems is that the constants $C_{N, (i_{1},\ldots,i_{N})} \tilde{\varphi}^{\mu \kappa}$ and the coefficients $Q_{j}$ in Theorem 9 have uniform bounds in $\sigma$.

As a consequence of Theorem 9 we obtain for arbitrary $\tilde{N} \in \mathbb{N}$

$$|J_{i_{1},\ldots,i_{N}}^\varphi(\nu) - (2\pi |\nu|^{-1})^\kappa Q_{0}(i_{1},\ldots,i_{N}) \tilde{\varphi}^{\mu \kappa} ; a_{i_{1},\ldots,i_{N}} \Phi_{i_{1},\ldots,i_{N}})|$$

$$
\leq \left|J_{i_{1},\ldots,i_{N}}^\varphi(\nu) - (2\pi |\nu|^{-1})^\kappa \sum_{l=0}^{\tilde{N}-1} |\nu|^{-l} Q_{l}(i_{1},\ldots,i_{N}) \tilde{\varphi}^{\mu \kappa} ; a_{i_{1},\ldots,i_{N}} \Phi_{i_{1},\ldots,i_{N}}) \right| + (2\pi |\nu|^{-1})^\kappa \sum_{l=1}^{\tilde{N}-1} |\nu|^{-l} |Q_{l}(i_{1},\ldots,i_{N}) \tilde{\varphi}^{\mu \kappa} ; a_{i_{1},\ldots,i_{N}} \Phi_{i_{1},\ldots,i_{N}})|$$

with constants $c_{i} > 0$ independent of both $\sigma$ and $\nu$. From this we deduce

$$\left|J_{i_{1},\ldots,i_{N}}^\varphi(\nu) - (2\pi / \mu)^\kappa \int_{0 < |\tau_{j}| < |\tau_{j}|^\kappa \in \sigma} Q_{0}^{N} |\tau_{j}|^{|c_j^2| + \sum_{l=1}^{j=1} d_{l}^{(j)}} - 1 - \kappa \log \det \delta_{i_{1},\ldots,i_{N}}(\sigma) \right| \left| \frac{d}{d\sigma} \right|$$

$$\leq c_{3} \mu^{-\tilde{N}} \int_{0 < |\tau_{j}| < |\tau_{j}|} Q_{0}^{N} |\tau_{j}|^{|c_j^2| + \sum_{l=1}^{j=1} d_{l}^{(j)}} - 1 - \kappa \log \det \delta_{i_{1},\ldots,i_{N}}(\sigma) \right| \left| \frac{d}{d\sigma} \right|$$

$$+ c_{4} \mu^{-\kappa} \sum_{l=1}^{\tilde{N}-1} |\nu|^{-l} \int_{0 < |\tau_{j}| < |\tau_{j}|} Q_{0}^{N} |\tau_{j}|^{|c_j^2| + \sum_{l=1}^{j=1} d_{l}^{(j)}} - 1 - \kappa \log \det \delta_{i_{1},\ldots,i_{N}}(\sigma) \right| \left| \frac{d}{d\sigma} \right|$$

$$\leq c_{5} \mu^{-\tilde{N}} \prod_{j=1}^{N} (- \log \varepsilon)^{q_{j}} \max \left\{ 1, \varepsilon^{c_j^2 + \sum_{l=1}^{j=1} d_{l}^{(j)}} \right\}$$

$$+ c_{6} \mu^{-\kappa} \prod_{l=1}^{\tilde{N}-1} (- \log \varepsilon)^{q_{l}} \max \left\{ 1, \varepsilon^{c_j^2 + \sum_{l=1}^{j=1} d_{l}^{(j)}} \right\},$$

where the exponents $q_{j}, q_{l}$ can take the values 0 or 1. Having in mind that we are interested in the case where $\mu \rightarrow +\infty$, we now set $\varepsilon = \mu^{-1/\tilde{N}}$. Taking into account Lemma 5 one infers that the right hand side of the last inequality can be estimated by a constant times

$$\mu^{-\kappa - 1} (\log \mu)^{N},$$

so that we finally obtain an asymptotic expansion for $I_{i_{1},\ldots,i_{N}}^\varphi(\nu) (\mu)$ by taking into account (44), and the fact that

$$(2\pi / \mu)^\kappa \int_{0 < |\tau_{j}| < |\tau_{j}|} Q_{0}^{N} |\tau_{j}|^{|c_j^2| + \sum_{l=1}^{j=1} d_{l}^{(j)}} - 1 - \kappa \log \det \delta_{i_{1},\ldots,i_{N}}(\sigma) \left| \frac{d}{d\sigma} \right|$$

$$= O(\mu^{-\kappa - 1}).$$

**Theorem 10.** Let the assumptions of the first main theorem be fulfilled. Then

$$I_{i_{1},\ldots,i_{N}}^\varphi(\nu) (\mu) = (2\pi / \mu)^\kappa E_{i_{1},\ldots,i_{N}}^\varphi(\nu) + O(\mu^{-\kappa - 1} (\log \mu)^{N}),$$

and the statement follows.
where the leading coefficient $\mathcal{L}^{\phi_1,\ldots,\phi_N}_{\chi_1,\ldots,\chi_N}$ is given by

$$
\mathcal{L}^{\phi_1,\ldots,\phi_N}_{\chi_1,\ldots,\chi_N} = \int_{\text{Crit}((\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}})} \frac{\alpha^{\phi_1,\ldots,\phi_N}_{\chi_1,\ldots,\chi_N} J^{\phi_1,\ldots,\phi_N}_{\chi_1,\ldots,\chi_N} \ d\text{Crit}((\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}})}{\det \text{Hess}((\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}})_{\mathcal{N}\text{Crit}((\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}})}}^{1/2}
$$

where $d\text{Crit}((\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}})$ denotes the induced Riemannian volume density.

\[\square\]

9. Statement of the main result

We can now state the main result of this paper. But before, we shall say a few words about the desingularization process. Consider the resolution of $\mathcal{N}$ constructed in Theorem 3, and denote the global morphism induced by the local transformations (44) by $Z: \tilde{X} \to X = T^* M \times G$. Consider further the local ideal $I_\Phi = (\Phi)$ generated by the phase function (10), together with the ideal sheaf $I_C \subset \mathcal{E}_X$ of (21). The derivative of $I_\Phi$ is given by $D(I_\Phi) = I_{C[T^* Y \times G]}$, while the implicit function theorem $\text{Sing} V_\Phi \subset V_\Phi \cap \text{Crit}(\Phi) = \text{Crit}(\Phi)$, where $V_\Phi$ denotes the vanishing set $V_\Phi$ of $\Phi$. The desingularization process carried out in Section 5 yields a partial monomialization of $I_\Phi$ according to the diagram

$$
Z^*(I_C) \cdot \mathcal{E}_{\tilde{x},\tilde{X}} \ni Z^*(I_\Phi) \cdot \mathcal{E}_{\tilde{x},\tilde{X}} = \prod_j \sigma_j^j \cdot Z_\ast^{-1}(I_\Phi) \cdot \mathcal{E}_{\tilde{x},\tilde{X}} \ni \prod_j \sigma_j^j \cdot (\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}}
$$

where $\tilde{x} \in \tilde{X}$. By Theorem 4, $D(Z_\ast^{-1}(I_\Phi))$ is a resolved ideal sheaf, and Theorem 5 shows that the weak transforms $(\chi_1,\ldots,\chi_N)\hat{\Phi}^{\text{wk}}$ have clean critical sets. This allowed us to derive asymptotics for the integrals $I^{\phi_1,\ldots,\phi_N}_{\chi_1,\ldots,\chi_N}(\mu)$ in Theorem 14. Nevertheless, it is easy to see that $Z_\ast^{-1}(I_\Phi)$ is not resolved. Furthermore, the inclusion (20) implies that $Z_\ast^{-1}(I_{C[T^* Y \times G]}) \subset D(Z_\ast^{-1}(I_\Phi))$. But since we do not have equality, this results only in a partial resolution $\tilde{C}$ of $C$. In particular, the induced global transform $Z: \tilde{C} \to C$ is in general not an isomorphism over the smooth locus of $C$. This is because of the fact that the centers of our monoidal transformations were only chosen over $M \times G$, to keep the phase analysis of the weak transform of $\Phi$ as simple as possible. In turn, the singularities of $C$ along the fibers of $T^* M$ were not completely resolved. Note that in order to obtain a partial monomialization of $I_\Phi$, we had to construct a strong resolution of $\mathcal{N}$ in $M = M \times G$, and not just a resolution of the $G$-action in $M$. As explained in Section 4, such a resolution always exists and is equivalent to a monomialization of the corresponding ideal sheaf. But in general, it would not be explicit enough to describe the asymptotic behavior of the integrals $I(\mu)$ introduced in (59). In particular, the so-called numerical data of $\zeta$ are not known a priori, which in our case are given in terms of the dimensions $c^{(i)}$ and $d^{(i)}$. This is the reason we were forced to construct an explicit resolution of $\mathcal{N}$, using as centers isotropy bundles over unions of maximally singular orbits.

Let us now return to our departing point, that is, the asymptotic behavior of the integrals $I(\mu)$, and the proof of Weyl’s law for the reduced spectral counting function $N_\zeta(\lambda)$. If $G$ acts on the chart $Y$ only with principal type $G/H_L$, we can directly apply the stationary phase theorem to obtain an expansion for $I(\mu)$. Let us therefore assume that this is not the case. We still have to
examine contributions to $I(\mu)$ coming from integrals of the form

$$I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu) = 
\int_{M_1(H_1) \times (-1,1)} \left[ \int_{\tilde{\mathcal{S}}(\mathcal{S}_1(\rho_{\mathcal{B}_1}))) \times (-1,1) \cdots \int_{\tilde{\mathcal{S}}(\mathcal{S}_1(\rho_{\mathcal{B}_N}))) \times (-1,1) \right] \cdots \frac{e^{i\mu \tau_1 \cdots \tau_N}}{a_{\mathcal{B}_1 \cdots \mathcal{B}_N}} \tilde{\mathcal{J}}_{\mathcal{B}_1 \cdots \mathcal{B}_N}(\mu),$$

where $\{\{H_1(\mu)\}, \ldots, (H_N(\mu))\}$ is an arbitrary totally ordered subset of non-principal isotropy types, $S_1(\rho_{\mathcal{B}_N})$ is the sphere of radius $\epsilon > 0$ in $\mathfrak{g}^\perp_{\mathcal{B}_N}$, while $a_{\mathcal{B}_1 \cdots \mathcal{B}_N}$ is an amplitude which is supposed to have compact support in a system of $(\mathcal{B}(\mathcal{B}_1), \ldots, \mathcal{B}(\mathcal{B}_N), \alpha(\mathcal{B}_1),\ldots)$-charts labeled by the indices $(\mathcal{A}_1, \ldots, \mathcal{A}_N)$, and

$$\tilde{\mathcal{J}}_{\mathcal{B}_1 \cdots \mathcal{B}_N}(\mu) = \prod_{i=1}^{\infty} |\tau_i| \left[ e^{(\gamma) + \sum_j d^{(\gamma)} - 1} \tilde{\mathcal{J}}_{\mathcal{B}_1 \cdots \mathcal{B}_N} \right],$$

$\tilde{\mathcal{J}}_{\mathcal{B}_1 \cdots \mathcal{B}_N}(\mu)$ being a smooth function which does not depend on the variables $\tau_i$. Now, a computation of the $\xi$-derivatives of $(\mathcal{B}(\mathcal{B}_1) \mathcal{B}(\mathcal{B}_1), \ldots)$ shows that $(\mathcal{B}(\mathcal{B}_1) \mathcal{B}(\mathcal{B}_1), \ldots)$-charts have no critical points there. Consequently, repeating the arguments of the previous section, and making use of the non-stationary phase theorem, see [28], Theorem 7.7.1, one computes for large $N \in \mathbb{N}$ that

$$|I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu)| \leq c_7 \mu^{-N} \int_{\epsilon < |\tau_i| < 1} \prod_{j=1}^{N} |\tau_j| \left[ e^{(\gamma) + \sum_j d^{(\gamma)} - 1} \tilde{\mathcal{J}}_{\mathcal{B}_1 \cdots \mathcal{B}_N} \right] + c_8 \mu^{-\kappa} \leq c_9 \max \{ \mu^{-N \kappa} \},$$

where we took $\epsilon = \mu^{-1/N}$. Choosing $N$ large enough, we therefore conclude that

$$|I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu)| = O(\mu^{-\kappa}).$$

As a consequence of this we see that, up to terms of order $O(\mu^{-\kappa-1})$, $I(\mu)$ can be written as a sum

$$I(\mu) = \sum_{N=1}^{\Lambda - 1} \sum_{\mathcal{B}_1 \cdots \mathcal{B}_N} I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu) + \sum_{N=1}^{\Lambda - 1} \sum_{\mathcal{B}_1 \cdots \mathcal{B}_N} I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu),$$

where the first term is a sum over maximal, totally ordered subsets of non-principal isotropy types, while the second term is a sum over totally ordered subsets of non-principal isotropy types. The asymptotic behavior of the integrals $I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu)$ has been determined in the previous section, and using Lemma 3, it is not difficult to see that the integrals $I_{\mathcal{B}_1 \cdots \mathcal{B}_N}^{\mathcal{A}_1 \cdots \mathcal{A}_N}(\mu)$ have analogous asymptotic descriptions. This leads us to the following

**Theorem 11.** Let $M$ be a connected, closed Riemannian manifold, and $G$ a compact, connected Lie group $G$ acting isometrically and effectively on $M$. Consider the oscillatory integral

$$I(\mu) = \int_{\mathcal{X}(\mathcal{X})} e^{i\mu \Phi(x,\xi,\eta)} a(x, x, \xi, \eta) d\mathcal{d}(\mathcal{V}^\ast)A(x, \xi), \quad \mu \to +\infty,$$

where $(\kappa, Y)$ are local coordinates on $M$, $d(\mathcal{V}^\ast)$ is the canonical volume density on $T^\ast Y$, and $d\mathcal{d}$ the volume density on $G$ with respect to some left invariant metric on $G$, while $a \in C_c^\infty(Y \times T^\ast Y \times G)$ is an amplitude, and $\Phi(x,\xi,\eta) = \kappa(x) - \kappa(gx, \xi).$ Then $I(\mu)$ has the asymptotic expansion

$$I(\mu) = (2\pi/\mu)^\kappa L_0 + O(\mu^{-\kappa-1}(\log \mu)^{\Lambda - 1}), \quad \mu \to +\infty.$$
Here $\kappa$ is the dimension of an orbit of principal type in $M$, $\Lambda$ the maximal number of elements of a totally ordered subset of the set of isotropy types, and the leading coefficient is given by

$$
(68) \quad L_0 = \int_{\text{Reg} C} \frac{a(gx, x, \xi, g)}{\det \Phi^\mu(x, \xi, g)|_{N(x, \xi, g)} \text{Reg} C}^{1/2} d(\text{Reg} C)(x, \xi, g),
$$

where $\text{Reg} C$ denotes the regular part of $C = \{(x, \xi, g) \in \Omega \times G : g \cdot (x, \xi) = (x, \xi)\}$, and $d(\text{Reg} C)$ the induced volume density. In particular, the integral over $\text{Reg} C$ exists.

**Remark 3.** Since $M$ is compact, $T^*M$ is a paracompact manifold, admitting a Riemannian metric. The restriction of the Riemannian metric on $T^*M \times G$ to $\text{Reg} C$ then induces a volume density $d(\text{Reg} C)$ on $\text{Reg} C$. Note that equation (68) in particular means that the obtained asymptotic expansion for $I(\mu)$ is independent of the explicit partial resolution we used. The amplitude $a(gx, x, \xi, g)$ might depend on $\mu$ as in the expression for $O(\mu^{-2})$ in Theorem 1. But as explained in the proof of Theorem 9, this has no influence on the final asymptotics.

**Proof.** Assume that $G$ acts on $Y$ with several orbit types. By Theorem 10 and (66) one has

$$
I(\mu) = (2\pi/\mu)^\lambda L_0 + O(\mu^{-\kappa - 1}(\log \mu)^{\Lambda - 1}), \quad \mu \to +\infty,
$$

where $L_0$ is given as a sum of integrals of the form (65), and similar expressions for the leading terms of the integrals $\int_{i_1, \ldots, i_{\lambda-1}} I(\mu)$. It therefore remains to show the equality (68). For this, let us introduce certain cut-off functions for the singular part $\text{Sing} \Omega$ of $\Omega$. Denote the Riemannian distance on $T^*M$ by $|\cdot|$, and let $K$ be a compact subset in $T^*M$, $\varepsilon > 0$. Then we define

$$
(S\text{ing} \Omega \cap K)_\varepsilon = \{\eta \in T^*M : |\eta - \eta'| < \varepsilon \text{ for some } \eta' \in \text{Sing} \Omega \cap K\}.
$$

By using a partition of unity, one can show the existence of a test function $u_\varepsilon \in C^\infty((\text{Sing} \Omega \cap K)_\varepsilon)$ satisfying $u_\varepsilon = 1$ on $(\text{Sing} \Omega \cap K)_\varepsilon$, see [25], Theorem 1.4.1. We then have the following

**Lemma 6.** Let $a \in C^\infty(Y \times T^*Y \times G)$, $K$ be a compact subset in $T^*M$ such that $\text{supp}(x, \xi) a \subset K$, and $u_\varepsilon$ as above. Then the limit

$$
(69) \quad \lim_{\varepsilon \to 0} \frac{\int_{\text{Reg} C} a(gx, x, \xi, g)(1-u_\varepsilon)(x, \xi)}{|\det \Phi^\mu(x, \xi, g)|_{N(x, \xi, g)} \text{Reg} C}^{1/2} d(\text{Reg} C)(x, \xi, g)
$$

exists and is equal to $L_0$.

**Proof of Lemma 6.** We define

$$
I_\varepsilon(\mu) = \int_{T^*Y} \int_G e^{i\mu \Phi(x, \xi, g)} a(gx, x, \xi, g)(1 - u_\varepsilon)(x, \xi) \, dg \, d(\text{Reg} C)(x, \xi).
$$

Since $(x, \xi, g) \in \text{Reg} C$ implies $(x, \xi) \in \text{Sing} \Omega$, a direct application of the generalized stationary phase theorem for fixed $\varepsilon > 0$ gives

$$
(70) \quad |I_\varepsilon(\mu) - (2\pi/\mu)^\lambda L_0(\varepsilon)| \leq C \varepsilon \mu^{-\kappa - 1},
$$

where $C \varepsilon > 0$ is a constant depending only on $\varepsilon$, and

$$
L_0(\varepsilon) = \int_{\text{Reg} C} \frac{a(gx, x, \xi, g)(1 - u_\varepsilon)(x, \xi)}{|\det \Phi^\mu(x, \xi, g)|_{N(x, \xi, g)} \text{Reg} C}^{1/2} d(\text{Reg} C)(x, \xi, g).
$$

On the other hand, applying our previous considerations to $I_\varepsilon(\mu)$ instead of $I(\mu)$, we obtain again an asymptotic expansion of the form (70) for $I_\varepsilon(\mu)$, with $\mu^{-\kappa - 1}(\log \mu)^{\Lambda - 1}$ instead of $\mu^{-\kappa - 1}$, where now the first coefficient is given by a sum of integrals of the form (65) with $a$ replaced by $a(1 - u_\varepsilon)$. Since the first term in the asymptotic expansion (70) is uniquely determined, the two expressions for $L_0(\varepsilon)$ must be identical. The statement of the lemma now follows by the Lebesgue theorem on bounded convergence. \[\square\]
Remark 4. Note that the existence of the limit in (69) has been established by partially resolving the singularities of the set $\mathcal{C}$, the corresponding limit being given by $L_0$.

End of proof of Theorem 1. Let now $a^+ \in C_c^\infty(Y \times T^*X \times G, \mathbb{R}^+)$. Since one can assume that $|u_\varepsilon| \leq 1$, the lemma of Fatou implies that

$$
\int_{\text{Reg } \mathcal{C}} \lim_{\varepsilon \to 0} \frac{a^+(gx, x, \xi, g)(1 - u_\varepsilon)(x, \xi)}{\det \Phi'(x, \xi, g)|_{N(x, \xi, g)_{\text{Reg } \mathcal{C}}}} d(\text{Reg } \mathcal{C})(x, \xi, g)
$$

is majorized by the limit in (69), with $a$ replaced by $a^+$. Lemma 5 then implies that

$$
\int_{\text{Reg } \mathcal{C}} \frac{a^+(gx, x, \xi, g)}{\det \Phi'(x, \xi, g)|_{N(x, \xi, g)_{\text{Reg } \mathcal{C}}}} d(\text{Reg } \mathcal{C})(x, \xi, g) < \infty.
$$

Choosing now $a^+$ to be equal 1 on a neighborhood of the support of $a$, and applying the theorem of Lebesgue on bounded convergence to the limit (69), we obtain equation (68).

We shall now state the main result of this paper.

Theorem 12. Let $M$ be a compact, connected, $n$-dimensional Riemannian manifold without boundary, and $G$ a compact, connected Lie group, acting effectively and isometrically on $M$. Let further

$$
P_0 : C^\infty(M) \to L^2(M)
$$

be an invariant, elliptic, classical pseudodifferential operator of order $m$ on $M$ with principal symbol $p(x, \xi)$, and assume that $P_0$ is positive and symmetric. Denote by $P$ its unique self-adjoint extension and set

$$
N_\chi(\lambda) = d_\chi \sum_{t \in \Lambda} \text{mult}_{\chi}(t),
$$

where $\text{mult}_{\chi}(t)$ stands for the multiplicity of the unitary irreducible representation $\pi_{\chi}$ corresponding to the character $\chi \in \hat{G}$ in the eigenspace $E_t$ of $P$ belonging to the eigenvalue $t$. Let $\kappa$ be the dimension of a $G$-orbit of principal type, $\Lambda$ the maximal number of elements of a totally ordered subset of the set of isotropy types, and assume that $n - \kappa \geq 1$. Then

$$
N_\chi(\lambda) = \frac{d_\chi[\pi_{\chi}|H : 1]}{(2\pi)^{n-\kappa} \text{vol } [(\Omega \cap S^*M)/G]} \lambda^{n-\kappa} \Lambda_{\kappa} + O(\lambda^{n-\kappa-1}(\log \lambda)^{\Lambda_{\kappa}}), \quad \lambda \to +\infty,
$$

where $d_\chi$ is the dimension of the irreducible representation $\pi_{\chi}$, $[\pi_{\chi}|H : 1]$ the multiplicity of the trivial representation in the restriction of $\pi_{\chi}$ to a principal isotropy group $H$, and $S^*M = \{(x, \xi) \in T^*M : p(x, \xi) = 1\}$, while $\omega = \mathbb{J}^{-1}(0)$ is the zero level of the momentum map $\mathbb{J} : T^*M \to g^*$ of the underlying Hamiltonian action.

Proof. Let $g \in C_c^\infty(-\delta, \delta)$ and $\delta > 0$ be sufficiently small. Theorems 11 and 11 together yield

$$
\hat{\sigma}_\chi(\varphi e^{t(\cdot)n}) = d_\chi \hat{\sigma}_\chi(\varphi)\mathcal{L}(\mu/2\pi)^{n-\kappa-1} + O(\mu^{n-\kappa-2}(\log \mu)^{\Lambda_{\kappa}-1}),
$$

where

$$
\mathcal{L} = \lim_{\varepsilon \to 0} \sum_{\gamma} \int_{\text{Reg } \mathcal{C}} \frac{u_{\gamma, \varepsilon}(x, \xi, g)}{\det \Phi'(x, \xi, g)|_{N(x, \xi, g)_{\text{Reg } \mathcal{C}}}} d(\text{Reg } \mathcal{C})(x, \xi, g),
$$

and $u_{\gamma, \varepsilon}(x, \xi, g) = \chi(g)f_\varepsilon(x, \xi, g)$. In order to compute $\mathcal{L}$, let us note that for any smooth, compactly supported function $\alpha$ on $\Omega \cap T^*Y$, one has the formula

$$
\int_{\text{Reg } \mathcal{C}} \frac{\chi(g)\alpha(x, \xi, g)}{\det \Phi'(x, \xi, g)|_{N(x, \xi, g)_{\text{Reg } \mathcal{C}}}} d(\text{Reg } \mathcal{C})(x, \xi, g) = [\pi_{\chi}|H : 1] \int_{\text{Reg } \Omega} \alpha(x, \xi, g) \frac{d\text{Reg } \Omega}{\text{vol } O(x, \xi)}.
$$

\[4\text{If } n - \kappa \geq 2, the error term is slightly better, namely } O(\lambda^{n-\kappa-1}(\log \lambda)^{\Lambda_{\kappa}-1}).\]
Thus, let \( \hat{\nu} \) denote the equivariant spectral counting function of \( N \). Here we took into account that by Proposition \( \[2 \] ) the set \( \{ (x, \xi) \in \text{Reg} \Omega : x \in \text{Sing} M \} \) has measure zero with respect to the induced volume form on \( \text{Reg} \Omega \), compare \([12], \text{Lemma 3} \). Next, let

\[
N^Q_X(\mu) = dx \sum_{\nu \leq \mu} \mu_X(\nu) = \sum_{\nu \leq \mu} m^Q_X(\mu_j), \quad m^Q_X(\mu_j) = d_X \mu_X(\mu_j)/\dim E^Q_{\mu_j},
\]

denote the equivariant spectral counting function of \( Q = P^{1/n} \). An asymptotic description for \( N^Q_X(\mu) \) can then be deduced from the one of \( \hat{\sigma}_X(\hat{\nu}e^{i(\cdot)\mu}) \) by a classical Tauberian argument \([2]\). Thus, let \( g \in C_c^\infty((-\delta, \delta)) \) be such that \( 1 = \int g(s) \, ds = 2\pi g(0) \). Then

\[
N^Q_X(\mu) = \int_{-\infty}^{+\infty} N^Q_X(\mu-s) \hat{g}(s) \, ds + \int_{-\infty}^{+\infty} [N^Q_X(\mu) - N^Q_X(\mu-s)] \hat{g}(s) \, ds =: H_X(\mu) + R_X(\mu).
\]

\[
H_X(\mu) = \int_{-\infty}^{+\infty} N^Q_X(s) \hat{g}(\mu-s) \, ds
\]

is a \( C^\infty \)-function, and by expressing its derivative by a Stieltjes integral one obtains

\[
\frac{dH_X}{d\mu}(\mu) = \int_{-\infty}^{+\infty} \frac{\partial}{\partial \mu} \hat{g}(\mu-s) N^Q_X(s) \, ds = -\int_{-\infty}^{+\infty} \frac{\partial}{\partial s} \hat{g}(\mu-s) N^Q_X(s) \, ds
\]

\[
= \int_{-\infty}^{+\infty} \hat{g}(\mu-s) dN^Q_X(s) = \sum_{j=1}^{\infty} m^Q_X(\mu_j) \hat{g}(\mu-\mu_j) = \hat{\sigma}_X(\hat{\nu}e^{i(\cdot)\mu}),
\]

where we took into account that, since \( \sigma_X \in S'(\mathbb{R}) \), \( N^Q_X(\mu) \) is polynomially bounded, and \( \hat{g}(s) = \hat{g}(-s) \). Now, in addition, let \( \hat{g} \) be such that \( \hat{g} \geq 0 \), and \( \hat{g}(s) \geq c_{10} > 0 \) for \( |s| \leq 1 \). Then, for \( \mu \in \mathbb{R} \),

\[
N^Q_X(\mu+1) - N^Q_X(\mu) \leq \sum_{|\mu-\mu_j| \leq 1} m^Q_X(\mu_j) \leq \frac{1}{c_{10}} \sum_{j=1}^{\infty} m^Q_X(\mu_j) \hat{g}(\mu-\mu_j).
\]

From \( \hat{\sigma}_X(\hat{\nu}e^{i(\cdot)\mu}) = O(\mu^{n-\kappa-1}) \) one then infers that \( R_X(\mu) = O(\mu^{n-\kappa-1}) \) as \( 1 \leq \mu \to +\infty \). On the other hand, since \( \hat{\sigma}_X(\hat{\nu}e^{i(\cdot)\mu}) \) is rapidly decaying as \( \mu \to -\infty \), integration gives

\[
H_X(\mu) = \int_1^{\mu} \hat{\sigma}_X(\hat{\nu}e^{i(\cdot)s}) \, ds + c_{11} = \frac{d \mathcal{L}}{n-\kappa} (\mu/2\pi)^{n-\kappa} + \begin{cases} O(\mu^{n-\kappa-1}(\log \mu)^{\lambda-1}), & n-\kappa \geq 2, \\ O((\log \mu)^{\lambda}), & n-\kappa = 1, \end{cases}
\]

as \( 1 \leq \mu \to +\infty \), while \( R_X(\mu) \), \( H_X(\mu) \to 0 \) as \( \mu \to -\infty \). The proof of the theorem is now complete, since by the spectral theorem, \( N_X(\lambda) = N^Q_X(\lambda^{1/n}) \). \( \square \)
10. On the Spectrum of $\Gamma \setminus G$

As an application, we shall consider the spectrum of a discrete, uniform subgroup $\Gamma$ of a connected, semisimple Lie group $G$ with finite center. Thus, let $\theta$ be a Cartan involution of $G$, and $g = \mathfrak{t} \oplus \mathfrak{p}$ the decomposition of the Lie algebra $\mathfrak{g}$ of $G$ into the eigenspaces of $\theta$. Let $K$ be the analytic subgroup corresponding to $\mathfrak{t}$, which is a maximal compact subgroup of $G$. Since $\Gamma$ is a uniform lattice, $M = \Gamma \setminus G$ is a closed manifold. By definition, $\theta$ is an involutive automorphism of $\mathfrak{g}$ such that the bilinear form $\langle X, Y \rangle_\theta = -\langle X, \theta Y \rangle$ is strictly positive definite, where $\langle X, Y \rangle = \text{tr}(\operatorname{ad} X \circ Y)$ denotes the Killing form on $\mathfrak{g}$. The form $\langle \cdot, \cdot \rangle_\theta$ defines a left-invariant metric on $G$, and by requiring that the the projection $G \to M$ is a Riemannian submersion, we obtain a Riemannian structure on $M$. Since $\operatorname{Ad}(K)$ commutes with $\theta$, and leaves invariant the Killing form, $K$ acts on $G$ and on $M$ from the right in an isometric and effective way. Note that the isotropy group of a point $Tg \in M$ is conjugate to the finite group $gKg^{-1} \cap \Gamma$. Hence, all $K$-orbits in $M$ are either principal or exceptional. Since the maximal compact subgroups of $G$ are precisely the conjugates of $K$, exceptional $K$-orbits arise from elements in $\Gamma$ of finite order. If $\Gamma$ is torsion-free, meaning that no non-trivial element $\gamma \in \Gamma$ is conjugate in $G$ to an element of $K$, there are no exceptional orbits. In this case the action of $\Gamma$ on $G/K$ is free, and $G/K$ becomes a smooth manifold of dimension $n - d$, where $n = \dim M$, and $d = \dim K$. As an immediate consequence of Theorem 12 we now obtain

**Corollary 2.** Let $P_0 : C^\infty(\Gamma \setminus G) \to L^2(\Gamma \setminus G)$ be a $K$-invariant, elliptic, classical pseudodifferential operator of order $m$ on $\Gamma \setminus G$, and assume that $P_0$ is positive and symmetric. Denote by $P$ its unique self-adjoint extension, and let $N_\chi(\lambda)$ be the reduced spectral counting function of $P$. Then, for each $\chi \in \hat{K}$,

$$N_\chi(\lambda) = \frac{d_{\chi}(\pi_\chi)|\mu; 1|}{(n-d)(2\pi)^{n-d}} \text{vol}([\Omega \cap S^* (\Gamma \setminus G))/K] \lambda^{\frac{n-d}{2}} + O\left(\lambda^{\frac{n-d}{2}(\log \lambda)^{\Lambda-1}}\right), \quad \lambda \to +\infty,$$

where $H \subset K$ is a principal isotropy group of the $K$-action on $\Gamma \setminus G$, and $\Lambda$ is bounded by the number of $\Gamma$-conjugacy classes of elements of finite order in $\Gamma$.

Under the assumption that $\Gamma$ has no torsion, this result was derived previously by Duistermaat–Kolk–Varadarajan [15] for the Laplace–Beltrami operator $\Delta$ on $L^2(\Gamma \setminus G/K) \simeq L^2(\Gamma \setminus G)^K$, i.e. in case that $\chi$ corresponds to the trivial representation. They proved this by studying the spectrum on $L^2(\Gamma \setminus G/K)$ of the whole algebra $\mathcal{D}(G/K)$ of $G$-invariant differential operators on $G/K$, which is defined as follows. Let $G = KAN$ be an Iwasawa decomposition of $G$, $a$ the Lie algebra of $A$, and $W$ the Weyl group of $(\mathfrak{g}, \mathfrak{a})$. Since $\mathcal{D}(G/K)$ is commutative, there is an orthogonal decomposition of $L^2(\Gamma \setminus G/K)$ into finite dimensional subspaces of smooth simultaneous eigenfunctions of $\mathcal{D}(G/K)$. Now, each homomorphism from $\mathcal{D}(G/K)$ to $\mathbb{C}$ is precisely of the form $\chi_\mu : \mathcal{D}(G/K) \to \mathbb{C}$, where $\mu \in a_\mathfrak{C}/W$. The spectrum $\Lambda(\Gamma)$ of $\mathcal{D}(G/K)$ on $\Gamma \setminus G/K$ is then defined as the set of all $\mu \in a_\mathfrak{C}/W$ for which there exists a non-zero $\varphi \in C^\infty(\Gamma \setminus G/K)$ with $D\varphi = \chi_\mu(D)\varphi$ for all $D \in \mathcal{D}(G/K)$. The main result of [15] is a description of the asymptotic growth of the tempered spectrum $\Lambda(\Gamma)_{\text{temp}} = \Lambda(\Gamma) \cap i\mathfrak{a}^*$, together with an estimate for the complementary spectrum $\Lambda(\Gamma) \setminus \Lambda(\Gamma)_{\text{temp}}$, using the Selberg trace formula, and the Paley–Wiener theorems of Gangolli and Harish-Chandra. From this, Weyl's law for $\Delta$ on $L^2(\Gamma \setminus G/K)$ follows readily, since the eigenvalue of $\Delta$ corresponding to $\mu \in \Lambda_{\text{temp}}(\Gamma)$ is essentially given by $||\mu||^2$.

Let now $P_0 : C^\infty(\Gamma \setminus G) \to L^2(\Gamma \setminus G)$ satisfy the conditions of Corollary 2, and in addition assume that it commutes with the right regular representation $R$ of $G$ on $L^2(\Gamma \setminus G)$. Then each eigenspace of $P$ becomes a unitary $G$-module. Since $\Gamma \setminus G$ is compact, $R$ decomposes into a direct
sum of irreducible representations of $G$ according to
\[ L^2(\Gamma \setminus G) \simeq \bigoplus_{\omega \in \hat{G}} m_\omega \mathcal{H}_\omega, \]
where $\hat{G}$ denotes the unitary dual of $G$, and $m_\omega$ the multiplicity of $\omega$ in $R$. In the same way, each eigenspace of $P$ decomposes into a direct sum of irreducible $G$-representations. Let $\operatorname{mult}_\omega(t)$ be the multiplicity of $\omega \in \hat{G}$ in the eigenspace $E_t$ of $P$ belonging to the eigenvalue $t$, and $[\omega|_K : \chi]$ the multiplicity of $\chi \in K$ in the $K$-representation obtained by restricting $\omega$ to $K$. Then
\[ \operatorname{mult}_\chi(t) = \sum_{\omega \in \hat{G}} \operatorname{mult}_\omega(t) [\omega|_K : \chi]. \]
Thus, the study of the reduced spectral counting function $N_\chi(\lambda)$ amounts to a description of the asymptotic multiplicities of those irreducible $G$-representations $\omega \in \hat{G}$ containing a certain $K$-type $\chi \in K$. As a consequence of Corollary 2 one now deduces

**Theorem 13.** Let $P_0 : C^\infty(\Gamma \setminus G) \to L^2(\Gamma \setminus G)$ be a $G$-invariant, elliptic, classical pseudodifferential operator of order $m$ on $\Gamma \setminus G$, and assume that $P_0$ is positive and symmetric. Denote by $\operatorname{mult}_\omega(t)$ the multiplicity of $\omega \in \hat{G}$ in the eigenspace $E_t$ belonging to the eigenvalue $t$ of the self-adjoint extension $P$ of $P_0$. Then, for each $\chi \in K$,
\[ \sum_{t \leq \lambda, \omega \in \hat{G}} \operatorname{mult}_\omega(t) [\omega|_K : \chi] = \frac{[\pi\chi|_H : 1]}{(n-d)(2\pi)^{n-d}} \operatorname{vol}[(\Omega \cap S^*(\Gamma \setminus G))/K] \lambda^{\frac{n-d}{2}} \]
\[ + O(\lambda^{\frac{n-d-1}{2}} (\log \lambda)^{\lambda^{-1}}), \quad \lambda \to +\infty, \]
where $n = \dim \Gamma \setminus G$, $d = \dim K$, and $\lambda$ is bounded by the number of $\Gamma$-conjugacy classes of elements of finite order in $\Gamma$.

\[ \square \]

**References**

[1] P. Albin and R. Melrose, *Equivariant cohomology and resolution*, arXiv Preprint 0907.3211, 2009.
[2] V. I. Arnold, S. M. Gusein-Zade, and A. N. Varchenko, *Singularities of differentiable maps*, vol. I, II, Birkhäuser, Boston, Basel, Berlin, 1988.
[3] M. F. Atiyah, *Resolution of singularities and division of distributions*, Comm. Pure Appl. Math. 23 (1970), 145–150.
[4] V.G. Avacumović, *Über die Eigenfunktionen auf geschlossenen Riemannschen Mannigfaltigkeiten*, Math. Z. 65 (1956), 327–344.
[5] I. N. Bernshtein, *Modules over a ring of differential operators. Study of the fundamental solutions of equations with constant coefficients*, Funkt. Anal. and its Appl. 5 (1971), no. 2, 89–101.
[6] I.N. Bernshtein and S. I. Gel’fand, *The meromorphic behavior of the function $P^\lambda$*, Funkt. Anal. and its Appl. 3 (1969), no. 1, 84–85.
[7] R. Bott, *On the iteration of closed geodesics and the Sturm intersection theory*, Comm. Pure Appl. Math. 9 (1956), 171–206.
[8] G.E. Bredon, *Introduction to compact transformation groups*, Academic Press, New York, 1972, Pure and Applied Mathematics, Vol. 46.
[9] J. Brüning, *Zur Eigenwertverteilung invarianter elliptischer Operatoren*, J. Reine Angew. Math. 339 (1983), 82–96.
[10] J. Brüning and E. Heintze, *Representations of compact Lie groups and elliptic operators*, Inv. math. 50 (1979), 169–203.
[11] T. Carleman, *Propriétés asymptotiques des fonctions fondamentales des membranes vibrantes*, C. R. Sém. Cong. Math. Scand. Stockholm, 1934, Lund, 1935, pp. 31–44.
[12] R. Cassanas and P. Ramacher, *Reduced Weyl asymptotics for pseudodifferential operators on bounded domains II. The compact group case*, J. Funct. Anal. 256 (2009), 91–128.
[13] M. Combescure, J.Ralston, and D. Robert, A proof of the Gutzwiller semiclassical trace formula using coherent states decomposition, Comm. Math. Phys. 202 (1999), 463–480.
[14] H. Donnelly, G-spaces, the asymptotic splitting of $L^2(M)$ into irreducibles, Math. Ann. 237 (1978), 23–40.
[15] J.J. Duistermaat, Oscillatory integrals, Lagrange immersions and unfolding of singularities, Comm. Pure Appl. Math. 27 (1974), 207–281.
[16] J.J. Duistermaat and V.W. Guillemin, The spectrum of positive elliptic operators and periodic bicharacteristics, Inv. Math. 29 (1975), no. 3, 39–79.
[17] J.J. Duistermaat and J.A. Kolk, Lie groups, Springer–Verlag, Berlin, Heidelberg, New York, 1999.
[18] J.J. Duistermaat, J. A. C. Kolk, and V. S. Varadarajan, Spectra of compact locally symmetric manifolds of negative curvature, Inv. Math. 52 (1979), 27–93.
[19] Z. El Houakmi and B. Helffer, Comportement semi-classique en présence de symétries I. Action d’un groupe de Lie compact, Asymptotic Anal. 5 (1991), no. 2, 91–113.
[20] L. Gårding, On the asymptotic distribution of eigenvalues and eigenfunctions of elliptic differential operators, Math. Scand. 1 (1953), 237–255.
[21] A. Grigis and J. Sjöstrand, Microlocal analysis for differential operators, London Mathematical Society Lecture Note Series, vol. 196, Cambridge University Press, 1994.
[22] V. Guillemin and A. Uribe, Reduction and the trace formula, J. Diff. Geom. 32 (1990), no. 2, 315–347.
[23] B. Helffer and D. Robert, Étude du spectre pour un opérateur globalement elliptique dont le symbole de Weyl présente des symétries I. Action des groupes fins, Amer. J. Math. 106 (1984), 1199–1236.
[24] ______, Étude du spectre pour un opérateur globalement elliptique dont le symbole de Weyl présente des symétries II. Action des groupes de Lie compacts, Amer. J. Math. 108 (1986), 973–1000.
[25] S. Helgason, Differential geometry, Lie groups, and symmetric spaces, American Mathematical Society, Providence Rhode Island, 2001.
[26] H. Hironaka, Resolution of singularities of an algebraic variety over a field of characteristic zero, Ann. of Math. 79 (1964), 109–326.
[27] L. Hörmander, The spectral function of an elliptic operator, Acta Math. 121 (1968), 193–218.
[28] L. Hörmander, The analysis of linear partial differential operators, vol. I, Springer–Verlag, Berlin, Heidelberg, New York, 1983.
[29] ______, The analysis of linear partial differential operators, vol. IV, Springer–Verlag, Berlin, Heidelberg, New York, 1985.
[30] K. Kawakubo, The theory of transformation groups, The Clarendon Press Oxford University Press, New York, 1991.
[31] J. Kollár, Lectures on resolution of singularities, Annals of Mathematical Studies, vol. 166, Princeton University Press, Princeton, New Jersey, 2007.
[32] B. Malgrange, Intégrales asymptotiques et monodromie, Ann. Sci. ENS 7 (1974), 405–430.
[33] E. Meinrenken and R. Sjamaar, Singular reduction and quantization, Topology 38 (1999), no. 4, 699–762.
[34] J.W. Milnor, Morse theory, Annals of Mathematical Studies, vol. 51, Princeton University Press, Princeton, New Jersey, 1963.
[35] S. Minakshisundaram and A. Pleijel, Some properties of the eigenfunctions of the Laplace operator on Riemannian manifolds, Canad. J. Math. 1 (1949), 242–256.
[36] J.P. Ortega and T.S. Ratiu, Momentum maps and Hamiltonian reduction, Progress in Mathematics, vol. 222, Birkhäuser Boston Inc., Boston, MA, 2004.
[37] P. Ramacher, Reduced Weyl asymptotics for pseudodifferential operators on bounded domains I. The finite group case, J. Funct. Anal. 255 (2008), 777–818.
[38] M.A. Shubin, Pseudodifferential operators and spectral theory, 2nd edition, Springer–Verlag, Berlin, Heidelberg, New York, 2001.
[39] R. Sjamaar and E. Lerman, Stratified symplectic spaces and reduction, Ann. of Math. 134 (1991), 375–422.
[40] S. Sternberg, Lectures on differential geometry, Prentice-Hall, Inc., Englewood Cliffs, New Jersey, 1965.
[41] V.S. Varadarajan, The method of stationary phase and applications to geometry and analysis on Lie groups, Algebraic and analytic methods in representation theory, Persp. Math., vol. 17, Academic Press, 1997, pp. 167–242.
[42] H. Weyl, Das asymptotische Verteilungsgesetz der Eigenwerte linearer partieller Differentialgleichungen (mit einer Anwendung auf die Theorie der Hohlraumstrahlung), Math. Ann. 71 (1912), 441–479.

PABLO RAMACHER, PHILIPPS-UNIVERSITÄT MARBURG, FACHBEREICH MATHEMATIK UND INFORMATIK, HANS-MEERWEIN-STR., 35032 MARBURG, GERMANY
E-mail address: ramacher@mathematik.uni-marburg.de