High temperature spectroscopy of ensembles of nitrogen vacancy centers in diamond

Mohammed Attrash, Oleg Shtempluck, and Eyal Buks
Andrew and Erna Viterbi Department of Electrical Engineering, Technion, Haifa 32000 Israel

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We study the spectroscopy of an ensemble of negatively charged nitrogen-vacancy (NV$^-$) centers in diamond at high temperatures between room temperature and 700 K under high vacuum conditions. Spin resonances are studied using optical detection of magnetic resonance (ODMR), and optical spectroscopy is employed to study radiative transitions. Upon increasing the temperature the intensity of radiative decay in visible and infra-red decreased. In addition, the ODMR resonance frequencies were decreased, and the phonon line emission shifted to higher wavelengths. Density functional theory calculation of the zero-field splitting parameter ($D$) revealed that thermal expansion is not enough to explain the shift in the ODMR frequencies. Fitting the measured intensity of photo-luminescence with the theoretical predictions of the Mott-Seitz model yields the value of 0.22 eV for the energy barrier associated with nonradiative decay.

I. INTRODUCTION

Negatively charged nitrogen vacancy (NV$^-$) center in diamond is considered as a candidate for quantum technologies due to its unique physical properties, such as long coherence time at room temperature. The technique of optical detection of magnetic resonance (ODMR) allows monitoring spin resonances [1–4], whereas radiative transitions can be probed using optical spectroscopy. Many researchers investigated the temperature dependence of luminescence of NV [5–9]. It was found that the magnetic resonance frequencies were decreased after increasing the temperature. The shift of the frequency lines is attributed to the temperature dependence of the NV$^-$ zero-field splitting parameter $D$. This dependency was explained by a combination of two mechanisms, thermal expansion [10] and electron-phonon interactions [8]. It is claimed that the non-radiative processes shorten the excited-state lifetime at high temperatures [5, 11, 12]. Beside shift in frequencies, the photo-luminescence (PL) intensity was found to decrease at high temperatures. Two models were proposed to explain PL intensity decay: Mott-Seitz [13, 14] and Schon-Klasens [15–17]. In the Mott-Seitz model, the PL quenching at high temperature is due to the enhancement of the non-radiative process. On the other hand, in the Schon-Klasens model, the PL decreases due to radiationless recombination of holes and electrons at non-radiative recombination centers. Despite the extensive studies of the NV at high temperatures, some remaining issues were not thoroughly investigated.

In this work we employ both ODMR in the microwave (MW) band, and optical spectroscopy in the optical band of 500 – 1500 nm, to study ensembles of NV$^-$ centers in a high vacuum (HV) system at different temperatures (300 – 550 K). In addition, we investigate the PL as a function of temperature and calculate the energy barrier according to Mott-Seitz model of ensembles of NV$^-$ centers. The comparison between our experimental results and theory yields partial agreement.

II. EXPERIMENTAL SETUP

Type Ib high pressure high temperature (HPHT) single crystal diamond with a nitrogen concentration lower than 200 ppm was laser-cut along the [110] plane, polished and irradiated with 2.8 MeV electrons at a dose of $8 \times 10^{18}$ cm$^{-2}$, and annealed at 900$^\circ$C for 2 hours. The estimated NV$^-$ concentration is $3.3 \times 10^{17}$ cm$^{-3}$ [20]. The diamond was cleaned in an acid mixture of Perchloric,
Sulfuric, and Fuming Nitric acid for 1 hour.

The vacuum setup is schematically shown in Fig. 1. The vacuum system, which has a base pressure of $1 \times 10^{-6}$ Torr, includes an ultra HV heater (HeatWave Labs Inc., model 101491) connected to a thermocouple (type K) and a temperature controller (model 101303), MW loop antenna (LA), and a bare multi-mode optical fiber pointing to the diamond to collect PL emitted from the diamond. Liquid nitrogen was flowed in a metal finger attached to the LA and the optical fiber to avoid annealing damage. The optical fiber is split (1x2 Step-Index Multimode Fiber Optic Couplers - Thorlabs) to 3 components, which are connected to: photo-diode (PD) (Thorlabs, PDA100A), visible spectrometer (VS) (Thorlabs, CCS175, 500 – 1000 nm), and infrared spectrometer (IRS) (Ibsen photonics, ROCK NIR 900 – 1700 nm). The relative optical power delivered to the PD, VS and IRS is 0.5, 0.25 and 0.25, respectively. The spectrometers allow monitoring both the triplet-triplet transition at wavelength of 637 nm and the singlet-singlet transition at 1042 nm [see the schematic energy level diagram shown in Fig. 2(a)].

The NV$^-$ centers were excited by a 532 nm green laser delivered via free space. The LA is connected to a radio frequency signal source (RFSG), and the MW amplitude was modulated by a 151 Hz sine wave. The PD is attached to a 600 nm long-pass filter, and its signal was de-modulated by a lock-in amplifier. To apply magnetic field on the diamond, a cylindrical neodymium magnet is positioned outside the vacuum chamber using a motorized stage. To reduce laser intensity fluctuations, we used proportional integral derivative (PID) controller (SIM 960 - Stanford Research Systems). A beam splitter in front of the laser together with a PD are used to generate the PID input signal. The longtime optical intensity stability is improved by more than a factor of 10 (compared with the stability obtained without the PID).

### III. ODMR

Ignoring the hyper-fine interaction, the NV$^-$ spin triplet ground state Hamiltonian $\mathcal{H}_\text{NV}$ is given by

$$\frac{\mathcal{H}_\text{NV}}{\hbar} = \frac{DS_z^2}{\hbar^2} - \frac{\gamma_e B \cdot \mathbf{S}}{\hbar} + \frac{E_{\text{NV}} (S_z^2 + S_x^2)}{2\hbar^2}, \quad (1)$$
where $D$ and $E_{\text{NV}}$ are the zero field splitting parameters, which equal to $2\pi \times 2.88$ GHz and $2\pi \times 10$ MHz at 3.6 K, respectively \[20\]. $\mathcal{S}$ is the total spin operator $\mathcal{S} = \mathcal{S}_x \hat{x} + \mathcal{S}_y \hat{y} + \mathcal{S}_z \hat{z}$, $\gamma_e = 2\pi \times 28.03$ GHz T$^{-1}$ is the electron spin gyromagnetic ratio, $\mathbf{B}$ is the magnetic field, and $\mathcal{S}_i = \mathcal{S}_x \pm i \mathcal{S}_y$. In this Hamiltonian we neglect also the on-axis electric field and some strain components which contribute to the $\mathcal{S}_y^2$ term and the off-axis electric field and other strain components which contribute to $\mathcal{S}_x$ and $\mathcal{S}_y$ terms \[21\]. Under continuous laser excitation, the NV$^-$ is polarized to the spin state $m_s = 0$, which has a brighter PL \[22\].

Figure 2(a) shows a schematic energy level diagram of an NV$^-$ center. The structure of nitrogen vacancy center in diamond is shown in Fig. 2(b). Measured ODMR of our sample as a function of magnetic field is shown in Fig. 2(c). The eight lines that are observed in this figure represent the resonances related to spin transitions from $m_s = \pm 1$ to $m_s = 0$ for the four different NV$^-$ center directions. The dashed lines represent the theoretically calculated frequencies of the same resonances, which are obtained by numerically diagonalizing the Hamiltonian $\mathcal{H}_{\text{NV}}$ \[11\]. Note that two pairs of resonances, which are represented by the pink dash lines, are nearly degenerate (note that their ODMR signal is nearly double the signal of the other NV directions). Figure 2(d) reveals the PL measurement of the HPHT diamond sample before and after electron irradiation and annealing. The measurement was recorded by Micro-Raman spectrometer (LABRAM HR, HORIBA, Jobin Yvon) using 532 nm wavelength laser at room temperature (laser power was lower than 10 mW). The HPHT diamond reveals a peak at 573 nm related to diamond Raman peak \[18\, 19\], and another low intensity peak at 637 nm related to the zero phonon line of NV$^-$ [see Fig. 2(a)] \[23\, 24\]. After electron irradiation and annealing process, the intensity of NV$^-$ zero phonon line increased, beside the PL intensity increase in the entire range of 650 – 850 nm. The intensity of the NV$^0$ zero phonon line is very low in the treated diamond sample. The reason for relatively low NV$^0$ peak intensity may be related to low laser power, as was found in Ref. \[25\].

IV. ZERO FIELD SPLITTING $D$

To numerically investigate the shift in the zero field splitting $D$, density functional theory (DFT) \[27\, 28\] was used as implemented in the QUANTUM ESPRESSO software \[29\]. The defect structure was modeled with $3 \times 3 \times 2$ and $2 \times 2 \times 2$ super-cell of diamond (144 and 64 atoms, respectively) and Γ point sampling of Brillouin zone. The calculation were performed with projector augmented-wave (PAW) method \[30\] and Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional \[31\] with a plane-wave cutoff of 300 Ry. To investigate the thermal expansion effect on $D$, the lattice length of the structure was chosen to suit the temperature as found in previous experiments \[32\, 33\], and the structure was relaxed at constant lattice length until the total force on the atoms was lower than $26 \times 10^{-3}$ eV Å$^{-1}$. The zero field splitting $D$ is believed to be related to spin-spin interactions of unpaired electrons \[34\]. Its tensor components $D_{ab}$ are given by \[35\] 

$$ D_{ab} = \frac{(\gamma_e \hbar)^2}{2S(2S-1)} \frac{\mu_0}{4\pi} \sum_{i<j} \chi_{ij} \langle \psi_{ij} | \frac{1}{r^3} \delta_{ab} - 3 r_a r_b | \psi_{ij} \rangle, \tag{2} $$
where \( a, b \in \{x, y, z\} \) are Cartesian indices, \( \mu_0 \) is the magnetic permeability of free space, the summation includes all pairs of occupied Kohn-Sham orbitals, \( \chi_{ij} = \pm 1 \) for parallel and anti-parallel spins, respectively, \( S \) is the spin number, and \( \psi_j(r, r') \) is \( 2 \times 2 \) determinant which can be written as \( \psi_j(r, r') = 2^{-1/2}[\psi_i(r)\psi_j(r') - \psi_j(r')\psi_i(r)] \). A Python package named PyZFS was used to calculate \( D_{ab} \).  

According to our DFT calculations, the \( D \) at zero temperature equals to 2.84 GHz for \( 3 \times 3 \times 3 \) super-cell (2.77 GHz for \( 2 \times 2 \times 2 \) super-cell) while the experimental value at 3.6 K equals to 2.88 GHz. For comparison to other DFT results, \( D \) was found in the range of [2.7, 3] GHz [37–39]. The shift of \( D \) at a given temperature, relative to its value at zero temperature, is denoted by \( \Delta D \). Some studies calculated the shift \( \Delta D \) according to the thermal expansion at temperatures below 300 K [8, 10]. To account for our experimental results, we calculate \( \Delta D \) based on the thermal expansion in the range from absolute zero to 500 K.

In order to investigate the temperature effect on the ODMR, the diamond was measured at different temperatures under constant magnetic field. The ODMR at temperatures below 520 K [see Fig. 3(a)] reveals resonance lines related to NV\(^-\). Upon increasing the temperature, the resonance frequencies and their intensity were decreased, and above 520 K the resonance lines could not be resolved. It is expected that the shift in resonance lines is due to the decrease of \( D \). The measured shift \( \Delta D \) in \( D \) [see Fig. 3(b)] at 500 K equals to \(-32\) MHz and the slope \( dD/dT \) is of order of \(-100\) kHz K\(^{-1}\), in agreement with the value reported in [8]. The normalized slope \( dD/(DdT) \) is \(-4.3 \times 10^{-5}\) K\(^{-1}\), which is higher than what was found in experiments carried out at lower temperatures, for which the value of \(-2.61 \times 10^{-5}\) K\(^{-1}\) has been measured [7, 8].

The experimental results of \( D \) at temperatures lower than 300 K were taken from other reports [7, 8]. As can be observed from simulation results, \( D \) decreases as the temperature increases. The shift in \( D \) at 500 K compare to its value at 300 K according to our calculations is \(-10\) MHz (\(-4\) MHz for the \( 2 \times 2 \times 2 \) super-cell). As can be seen from the plot in Fig. 3(b), the thermal expansion is not enough to explain the shift in \( D \). The unit cell length \( a \) of diamond at 298 K and 506 K equals to 3.5668 \( \text{Å} \) and 3.5680 \( \text{Å} \), respectively [32], therefore, \( da/(adT) \approx 1.6 \times 10^{-6}\) K\(^{-1}\) in this range. This value is too small to account for the above-mentioned experimental finding that \( dD/(DdT) \) is \(-4.3 \times 10^{-5}\) K\(^{-1}\) in the same range.

According to a first-principles calculation reported in [40], the shift in \( D \) at 500 K is \(-18\) MHz, while the shift in our data is \(-32\) MHz. Good agreement is found between our experimental results, and the ones reported in Ref. [3], in which the zero-field splitting parameter \( D \) is expressed as a power series of the temperature \( T \).

The ODMR signal intensity can be derived from the rate equation that governs the time evolution of the spin decay in the intermediate singlet states \(^1\text{A}_1 \rightarrow ^3\text{E}\), can be observed at temperatures below 450 K.
polarization (see appendix A of Ref. [20]). As can be seen from Fig. 3, ODMR signals could not be experimentally resolved above a temperature of about 520 K. On the other hand, PL signals were measured at higher temperatures up to about 700 K (see Fig. 6). This observation suggests that the dominant mechanism responsible for the observed ODMR signal drop with temperature [see Fig. 3(a)], is dephasing, which enhances spin transverse relaxation, and consequently disables ODMR above 520 K.

V. OPTICAL SPECTROSCOPY

The fluorescence emission spectra of NV\(^{-}\) (measured by the VS) at different temperatures is shown in Fig. 4. The fluorescence can be resolved at temperatures up to 520 K. The spectra is shifted to longer wavelengths at higher temperatures. The shift in wavelength at which the emission intensity is maximized, which is denoted by \(\lambda_{\text{max}}\), equals to \(hc/60\text{ meV at 500 K, where} h\) is the Planck constant and \(c\) is the speed of light, and thus \(d\lambda_{\text{max}}/\lambda_{\text{max}}\) \(\approx 3.1 \times 10^{-4}\text{ K}^{-1}\). Fitting of the phonon shift was performed according to \(\Delta E_\nu(T) = \frac{-2A}{\exp \left(\frac{\hbar \omega_0}{2k_B T}\right) - 1}\),

where \(A\) depends on the diamond dispersion curves details, \(k_B\) is the Boltzmann’s constant, and \(\omega_0\) is the zone-center phonon frequency. The fitting yields that \(\hbar \omega_0 = 173 \pm 10\text{ meV}\) and \(A = 246 \pm 34\text{ meV}\). Note that the zone-center optical phonon Raman energy of diamond equals to 165 meV (1332 cm\(^{-1}\)) [22].

The IR spectra of the NV\(^{-}\) ensemble (measured by the IRS) at different temperatures is shown in Fig. 5. The peak at 1042 nm, which is associated with radiation decay from the \(^1\text{A}_1\) to \(^\text{E}_1\) levels [see Fig. 2(a)], is observed at room temperatures. Upon heating the diamond, the 1042 nm peak intensity decreased and became unresolved at temperatures above 450 K. Temperature dependence of the \(^1\text{A}_1\) to \(^\text{E}_1\) transition wavelength has been observed in [8]. However, this dependency cannot be resolved in our measurements due to limited IRS resolution. Another feature can be observed at 1350 nm. We find that this feature disappears when a 808 nm long pass filter is added, i.e. it originates from second order diffraction (note that 1350/2 = 675). The singlet NV\(^{-}\) transition (from \(^1\text{A}_1\) to \(^\text{E}_1\) states) wavelength is found to be less sensitive to temperature compare to the triplet transition (from \(^3\text{A}_2\) to \(^3\text{E}\) states) wavelength.

Finally, we discuss the PL intensity and energy barrier for the nonradiative process, which is denoted by \(U_b\). The PL intensity \(I_{\text{PL}}\) as a function of the temperature \(T\) is shown in Fig. 6. In the Mott-Seitz model this dependency of \(I_{\text{PL}}\) on \(T\) is given by [see Eq. (2) of Ref. [43], and Eq. (5) of Ref. [14]]

\[
I_{\text{PL}} = \frac{I_0}{1 + C \exp \left(-\frac{U_b}{k_B T}\right)} ,
\]

where \(I_0\) is the zero temperature intensity, and \(C\) is a constant. Fitting the measured PL intensity \(I_{\text{PL}}\) with Eq. 4 yields \(U_b = 0.22 \pm 0.05\text{ eV and } C = 200\). Note that \(U_b\) was found in ref. [8] to be 0.48 eV. The deviation between this result [5], which was obtained with a single NV\(^{-}\) center, and our extracted value for the energetic barrier, can perhaps be attributed to impurities near the single NV center that was studied in Ref. [8]. The effect of nearby impurities on coherence lifetime of NV\(^{-}\) centers has been studied in [45].

VI. SUMMARY

In this study we investigated the optical properties of an ensemble of NV\(^{-}\) centers in diamond as a function of temperature. The NV\(^{-}\) emission in the visible \((\sim 700\text{ nm})\) and IR \((\sim 1042\text{ nm})\) regions decreased upon increasing the temperature, and the phonon line wavelength at maximum intensity of the NV\(^{-}\) \((650 - 800\text{ nm})\) increased. The \(D\) parameter decreased upon increasing the temperature, and the thermal expansion model is not enough to explain this behaviour. By applying the Mott-Seitz model we find that the energy barrier for nonradiative decay \(U_b\) is in the range \([0.17, 0.27]\) eV.

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