Modelling and thermal simulation of absorber-free quasi-simultaneous laser welding of transparent plastics

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Abstract
The growing demands on the quality of plastic components and the trend towards miniaturisation are posing a great challenge on plastics processing technology. As many complex components can no longer be manufactured in a single step, joining processes such as laser transmission welding are gaining in importance. In classic laser transmission welding, the joining partners have different optical properties. The upper joining partner is transparent in the laser wavelength range, whilst the lower partner is absorbent due to the addition of absorber materials. In medical and biotechnological applications, the addition of absorber materials is often undesirable due to strict biocompatibility requirements. If, on the other hand, radiation sources are used which emit radiation in the area of the natural absorption of the plastic (\(\lambda = 1600–2000\) nm), untreated transparent plastics can also be welded. In this work, a theoretical model will be presented to calculate the temperature distribution and progression during quasi-simultaneous welding using a thulium fibre laser (\(\lambda = 1940\) nm). A sensitivity analysis is carried out to investigate the influence of different parameters on the heat affected zone (HAZ). The simulated HAZ is then compared with the HAZ from the experimental work.

Keywords Plastic joining · Transparent polymers · Thermal simulation · Laser transmission welding

1 Introduction
The growing demands on the quality of plastic components and the trend towards miniaturisation are posing a great challenge on plastics processing technology. As many complex components can no longer be manufactured in a single step, joining processes such as laser transmission welding are gaining in importance [1]. For example, microfluidic components such as lab-on-a-chip devices combine the functionalities of analytical instruments in a small chip. The higher portability allows these tests to be performed on site. The challenge here is the media-tight encapsulation of the narrow and complex channel structures with channel diameters less than 300 \(\mu\)m.

One promising approach is laser transmission welding. The temporal and spatial defined and contactless energy input leads to a precise welding process at low thermal and mechanical stress [2, 3]. In classic laser transmission welding, beam sources are used which emit radiation in the transparent wavelength range of most polymers (\(\lambda = 800–1100\) nm). Therefore, the modification of the optical components is necessary to melt the joining partners. Generally, the joining partners are welded as an overlap joint. A clamping device enables an exact and reproducible positioning of the joining partners. Furthermore, a thermal contact is achieved by applying a joining pressure. During the welding process, the incident laser beam is transmitted by the upper transparent joining partner and absorbed at the surface of the lower absorbent joining partner. The absorbing properties of the lower joining partner are generated by using an absorber material such as carbon black. The electromagnetic irradiation energy is converted into heat which leads to melting of both joining partners in the joining zone. A welding seam is then generated as soon as the melt solidifies [3, 4]. One important factor for the achievement of a high and reproducible welding seam quality is the amount of absorber material. Here, an exact dosage is decisive as fluctuations of the absorber amount lead to an uneven absorption and therefore to defects of the welding seam [5]. On
top of that, the black colouring of the component is inadmissible in areas such as biotechnology or medical engineering [6]. Alternatively, it is possible to use near-infrared absorbers which do not or only slightly affect the colour of the component. As the price is significantly higher, this has a huge impact on high volumetric compounds [7].

By using new beam sources which emit radiation in the area of the polymeric absorption bands (λ = 1600–2000 nm), it is nowadays possible to join transparent polymers without the use of absorbers. The melting is achieved by the excitation of the first harmonic of various bonds such as methyl and methylene groups. Promising application fields apart from the encapsulation of microfluidics is the sealing of polymer films, e.g. food packaging or sensitive electronic components such as OLEDs [7, 8].

With the help of thermal simulations, the temporally and spatially resolved temperature distribution and progression can be determined inside the material. In case of classic laser transmission welding, many works are available based on a finite element method (FEM). For instance, Kreimeier calculated the residual stress inside the welding seam due to the thermal expansion and shrinkage [9]. Other works deal with the influence of absorber materials such as carbon black or indium tin oxide on the temperature field and the dimensions of the HAZ. In case of absorber-free quasi-simultaneous welding, a sensitivity analysis is investigated during the welding process. In order to determine the influencing factors on the HAZ, a sensitivity analysis is carried out. Furthermore, the simulated results will be compared with the HAZ of the real welding experiments.

2 Thermal simulation of absorber-free laser transmission welding

2.1 Fundamentals

The transient temperature field in plastics can be described by Fourier’s differential equation [1]:

\[ \rho(T) c_p(T) \left( \frac{\partial T}{\partial t} + \nabla \cdot \mathbf{q} \right) = \nabla \cdot (k \nabla T) + Q \]

where \( \rho \) is the density, \( c_p \) is the heat capacity, \( k \) is the heat conductivity, and \( Q \) is the heat source.

The heat source \( Q \) is determined by the absorbed laser beam intensity during propagation through the material. In case that no scattering occurs inside the material (such as in amorphous polymers), the heat source can be described by the absorbed laser intensity [9]:

\[ Q = \alpha \cdot I(x, y, z, t, \alpha) \]

where \( \alpha \) is the absorption coefficient. The intensity profile of an ideal laser beam has a Gaussian distribution. It can be described as:

\[ I(r, z) = I_0(z) e^{-2 \left( \frac{\pi r}{w(z)} \right)^2} \]

\[ I_0(z) \text{ is defined as:} \]

\[ I_0(z) = I_0 \left( \frac{w_0}{w(z)} \right)^2 \]

where \( w_0 \) is the focal beam diameter. The beam radius as a function of the z-coordinate is determined by:

\[ w(z) = w_0 \sqrt{1 + \left( \frac{z}{z_R} \right)^2} \]

where \( z_R \) is the Rayleigh length and \( w_0 \) is the focal diameter.

The intensity is defined as power per area. Consequently, the laser power \( P_L \) can be calculated by the integral over the radial intensity distribution.

\[ P_L = \int_{\varphi=0}^{\pi} \int_{r=0}^{r=e} I(r, z = 0) \cdot r \, dr \, d\varphi \]

Solved after \( I_0 \), the equation results in:

\[ I_0 = \frac{2 \cdot P_L}{\pi \cdot w_0^2} \]

This results in the following spatial intensity distribution:

\[ I(r, z) = \frac{2 \cdot P_L}{\pi \cdot w_0^2} \left( \frac{w_0}{w(z)} \right)^2 e^{-2(z/\pi z_R)^2} \]
By a transformation into Cartesian coordinates and by consideration of the movement of the laser beam, \( r^2 \) can be reformulated as follows:

\[
r^2 = (x-v*t)^2 + y^2
\]

Considering the intensity reduction according to the Lambert Beer law and the reflection at the boundary surface, the following equation results for the heat source:

\[
Q = \frac{(1-R) \cdot \alpha}{\pi \cdot w(z)^2} \cdot e^{-\alpha \cdot z} \cdot \frac{2 \cdot P_L}{(v \cdot \pi \cdot w(z)^2 + \alpha^2)^2}
\]

### 2.2 Modelling of the welding process

The thermal simulation was carried out using the FEM software Comsol Multiphysics® (Comsol Inc., Stockholm, Sweden). Here, a polycarbonate (PC) was modelled as a sample with a thickness of 2 mm. In order to calculate the absorbed laser intensity inside the material, the optical properties were measured using the UV/VIS/NIR spectrometer Lambda 1050 (Perkin Elmer Inc., Waltham, MA). At 1940 nm, the absorption is 15%, the reflection is 9%, and the transmission is 76%. The sample is enclosed by a glass plate and an aluminium plate, which form the clamping device (see Fig. 1). About 85% of the laser radiation passes through the glass plate and is absorbed by the sample. The remaining radiation is reflected back into the sample by the aluminium plate. Roughness has been neglected for all components.

Since in the ideal case the temperature distribution in the plastic is symmetrical, only half of the sample was modelled to simplify matters. This shortens the calculation time. The beam properties were measured using the scanning diagnostics system Focus Monitor (Primes GmbH, Pfungstadt, Germany) and implemented in the simulation model. The seam length is 16 mm. In this work, the HAZ is defined as the area where the glass transition temperature of PC \((T_G = 145 \, ^\circ C)\) is exceeded. The value was taken from the literature [17].

The simulations were performed for different welding parameters in order to investigate their influence on the HAZ. For a better comparability, the line energy \( E \) was used. It is defined as:

\[
E = \frac{P}{v \cdot n}
\]

where \( P \) is the laser power, \( v \) is the feed rate and \( n \) is the number of repeats. As described in Chapter 2.1, the heat conduction equation consists of a convective and a diffusive component. The Péclet number describes the ratio between the feed rate, the beam diameter and the thermal diffusivity. It is an important parameter for characterising the temperature development:

\[
P\acute{e} = \frac{v \cdot w_0}{a}
\]

where \( v \) is the feed rate, \( w_0 \) is the beam diameter and \( a \) is the thermal diffusivity. The convective energy flow is influenced by the feed rate, whilst the diffusive term is determined by the thermal diffusivity. In case of \( P\acute{e} > > 1 \), the convective heat flow dominates, so that the temperature distribution corresponds to the moving heat source. This is the case with high feed rates. For \( P\acute{e} << 1 \), the diffusive heat flux is dominant, which results in a wider temperature field in the material.

### 3 Simulation results

#### 3.1 Simulation of the heating performance

In a first step, the temperature development in the interface between the upper and lower joining partners was simulated.
The welding parameters selected were a constant line energy of 2.5 J/mm and a laser power of 22.5 W. The number of passes was increased from \( n = 1 \) to \( n = 150 \). In order to keep the line energy constant, the feed rate must be adjusted accordingly. For \( n = 1 \), the feed rate is therefore 9 mm/s, whilst for \( n = 150 \), it increases to 1350 mm/s.

\[
E = \frac{22.5 \text{ W}}{9 \frac{\text{mm}}{s}} \times 1 = \frac{22.5 \text{ W}}{1350 \frac{\text{mm}}{s}} \times 150 = 2.5 \text{ J/mm}
\]

For \( n = 1 \), the highest temperature in the material is achieved with \( T = 291 \, ^\circ\text{C} \) (see Fig. 2). The short but strong heating of the material and the subsequent cooling process can also be seen. Overall, contour welding is characterised by a short interaction time between the laser beam and the material. For higher \( n \), the welding process can be divided into \( n \) heating and cooling phases. A heating phase \( t_h \) is followed by a cooling phase \( t_c \), during which heat is transferred to the aluminium and glass plates via the outer surfaces of the sample. From \( n = 5 \), a lower maximum temperature is reached in the material due to the cooling phases. In contrast to contour welding, the interaction times are longer. The results are shown in Fig. 2.

From \( n = 100 \), the cooling phase decreases due to the high feed rates to such an extent that a single heating phase can be assumed (\( t_c \approx 0 \)). Consequently, a quasi-simultaneous heating of the entire seam segment takes place. This shifts the ratio of the heat flow, and the convective term can be neglected. This in turn leads to a decrease of the Péclet number.

### 3.2 Simulation of the temperature distribution

The following figure (Fig. 3) shows the temperature distribution along the cross section of the sample. The sample is located between \( z = 1 \) mm and \( z = 3 \) mm, and the glass and aluminium plate is located above and below it. For \( n = 1 \), the highest temperatures are reached inside the polymer, whereas the material temperature decreases with increasing number of passes. The glass transition temperature is therefore exceeded in a smaller range. Consequently, the reduction of the vertical expansion of the HAZ can be achieved by means of quasi-simultaneous welding. For all parameters, the maximum temperature is not located in the interface, but in the upper joining partner. The HAZ is thus shifted in the direction of the upper joining partner. The reason for this is the higher thermal conductivity of aluminium compared with glass. Whilst the aluminium plate is not heated significantly, there is an increase in temperature in the glass plate. Another reason is that the laser intensity is highest on the surface of the upper joining partner.

In case of the HAZ width, the highest maximum temperature in the plastic is achieved with a value of 292 °C for \( n = 1 \). It can be seen that the temperature gradient is significantly greater than for all other parameters. The glass transition temperature is also exceeded in a wider range. This means that the highest HAZ width is achieved during contour welding. As the number of passes increases, the maximum temperature in the polymer decreases. A reduction in the HAZ width is also achieved. However, the reduction occurs to a lesser extent compared with the HAZ height. The reason for this is that the heat in the joining zone is not dissipated via the outer surfaces. Instead, the heat flows off into the surrounding materials.
material, which also leads to heating here. Contour welding, on the other hand, heats the surrounding material only slightly.

### 3.3 Sensitivity analysis

The results presented above show the influence of the process parameters number of passes \( n \) and feed rate \( v \) on the temperature distribution and the HAZ. However, the influence of the following parameters is unknown on the generation of the HAZ:

- Absorption coefficient \( \alpha \)
- Thermal conductivity of the polymer \( \lambda_{PC} \)
- Thermal conductivity of the glass plate \( \lambda_{glass} \)
- Beam diameter \( w_0 \)
- Density and heat capacity \( \rho \cdot c_p \)

In order to quantify the influence, a sensitivity analysis is carried out according to the one-at-a-time technique (OAT). Only one input variable is varied, whilst the remaining parameters remain unchanged. In total, each parameter was varied by \( \pm 30\% \), and the effect on the HAZ was investigated. The simulations were carried out for \( P = 2.5 \) W, \( n = 20 \) and \( v = 180 \) mm/s, so that the line energy is \( E = 2.5 \) J/mm. The results are shown in Fig. 4.

In a first step, the thermal conductivity \( \lambda_{PC} \) was varied. The lower the value, the stronger the effect of heat accumulation becomes and the less heat is dissipated into the surrounding material and the outer surfaces of the joining partners. With a 30% reduction of \( \lambda_{PC} \), the HAZ width increases by 21% and the HAZ height by 23.5%. However, if \( \lambda_{PC} \) is reduced, more heat will flow into the environment during the welding process. In this case, the HAZ width is reduced by 32.6% and the height by 31.3%. The absorption coefficient \( \alpha \) determines the absorbed laser energy along the beam propagation direction.

If \( \alpha \) is reduced by 30%, the absorbed radiation is not sufficient to form a HAZ. As a result, the HAZ width and height are changed by \( -100\% \). An increase leads to an increase in the HAZ width and height of 48.8% and 25%, respectively.

Of all investigated parameters, the thermal conductivity of the glass plate has the least influence on the formation of the

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**Fig. 3** Vertical (left) and horizontal temperature distribution (right)

**Fig. 4** Change of the HAZ width (left) and HAZ height (right) after the sensitivity analysis
HAZ. With a reduction, the HAZ width increases by 4.3%, while the HAZ height increases by 0.86%. If the thermal conductivity of the glass plate is increased, the HAZ height decreases by 1.7%, but the HAZ width also increases by 2.1%. Due to the small amount of the change, this can be attributed to measurement errors.

In addition to the absorption coefficient, the product of the density and the heat capacity of PC also have a major influence on the HAZ. A reduction of $\rho \cdot c_p$ leads to an increase of 47.9% in the HAZ width and 19% in the HAZ height. The HAZ width decreases by 56.2% with an increase of $\rho \cdot c_p$, while the HAZ height decreases by 40.5%.

In contour welding, the HAZ dimensions are mainly determined by the focal diameter. In the theoretical consideration of quasi-simultaneous welding, the effect is reversed. In the case of a reduction of the focal diameter, the HAZ width increases by 10.6%, while the HAZ height increases by 24.3%. If the focal diameter is increased, the HAZ width and height decrease by 4.2% and 19.1%, respectively. In order to explain this effect, it is necessary to observe the intensity and temperature distribution perpendicular to the scanning direction.

Figure 5 on the left shows the intensity distribution in the joining zone. With a focal diameter of 85 $\mu$m, the maximum intensity is $1.85 \times 10^4$ W/cm$^2$. If the focal diameter is increased to 152 $\mu$m, the maximum intensity drops to $0.58 \times 10^4$ W/cm$^2$. The intensity thus drops by approx. 66%. The higher focal diameter also leads to a widening of the intensity distribution. As a result, the intensity at a distance of ± 0.29 mm is greater at a focal diameter of 152 $\mu$m than at 85 $\mu$m. A reduction in the HAZ dimensions with a larger focal diameter is due to the fact that the threshold intensity required for melting is exceeded in a smaller range. This is also shown by the temperature distribution in Fig. 5 on the right. At $w_0 = 152$ $\mu$m, the maximum temperature in the joining zone is 165 °C, whilst at $w_0 = 85$ $\mu$m, the temperature is 201 °C. Also, with a higher focal diameter, the glass transition temperature is exceeded in a smaller range.

However, it should be noted that this is only a theoretical consideration of the welding process. Whether the above-mentioned effect also occurs in real welding tests must be proven in future studies.

4 Experimental validation

4.1 Experimental setup

The experimental setup is shown in Fig. 6. As testing material, PC with a thickness of 1 mm was used (Makrolon®, Covestro AG, Leverkusen, Germany). The samples were placed in a
clamping device which consists of a glass and an aluminium plate. This enables a thermal contact as well as a reproducible positioning. The samples were then irradiated using a thulium fibre laser with a maximum output power of $P = 120$ W (IPG Laser GmbH, Burbach, Germany). The laser beam was moved using the scanning system IntelliSCAN III 14 (Scanlab GmbH, Puchheim, Germany) and focused using a telecentric f-theta lens ($f = 100$ mm).

The welding seams were investigated using a polarisation microscope Axio Imager A2.m (Carl Zeiss Microscopy GmbH, Jena, Germany). Here, thin sections were made of the samples perpendicular to the welding seams using a rotary microtome RM2255 (Leica Biosystems Nussloch GmbH, Nussloch, Germany). A summary of the laser system and the beam properties can be found in Table 1.

### 4.2 Comparison between experimental and simulational results

In the following, the HAZ dimensions determined by the simulation are compared with the values from the real welding tests. Figure 7 (left) shows the height of the HAZ as a function of the number of passes. As the number of passes and feed rate increases, the diffusive energy flow dominates, and consequently more heat is dissipated via the outer surfaces of the joining partners. In case of the HAZ height, a value of 1.76 mm was reached for $n = 5$. For $n = 150$, the HAZ height was reduced to 1 mm which leads to a reduction of 43.4%. In the simulation on the other hand, a reduction of 67.4% was achieved. On average, the deviation between the simulated and the experimental HAZ height is 12.8%.

A similar trend can be observed for the HAZ width. The HAZ width is 0.4 mm for $n = 5$, whilst at $n = 150$, the HAZ width is reduced to 0.3 mm. This leads to a reduction of 33.2%. The simulated results, however, show a reduction of 46.3%. It can also be seen that there is a high deviation between the simulated and experimental HAZ width with a value of 48.5%.

Additionally, the simulation model predicts a smaller HAZ height compared with the experimental results. On the other hand, the simulated HAZ width is significantly higher than the real HAZ width.

A possible reason for the deviations is that simplifications in the simulation model were assumed to reduce the computing time. For example, a perfect thermal contact between both joining partners was assumed, and the roughness was therefore neglected. An ideal reflection of the laser radiation at the glass plate was also assumed. Another reason can be the influence of the interaction time between laser radiation and plastic on the melting temperature.

In a previous work by Boglea, it was indicated that the melting temperature is shifted towards higher temperatures with shorter interaction times [18]. For example, the melting temperature of a polymer is usually obtained by differential scanning calorimetry or thermogravimetric analysis measurements. Here, the heating rates are only 10–30 °C/min. In laser transmission welding, the heating rates can reach up to 150 °C/ms. In this work on the other hand, a possible increase

### Table 1 Summary of the laser system and the beam properties

| Laser system          | Beam properties |
|-----------------------|-----------------|
| Beam source           | IPG TLR-120-1940 |
| Scanner               | SCANLAB IntelliSCAN III 14 |
| Focal length          | 100 mm          |
| Focal diameter        | 116 μm          |
| Rayleigh length       | 1.97 mm         |
| Angle of divergence   | 1.65 °          |

![Fig. 7](image_url) Comparison between the simulation and the experimental results
of the melting temperature due to shorter interaction times was not considered.

5 Conclusion

In the present work, a model for the calculation of the temperature distribution during absorber-free quasi-simultaneous welding of plastics with a wavelength of 1940 nm was presented. It could be shown that the diffusive heat transport dominates in quasi-simultaneous welding. With higher feed rate and number of passes, the heat dissipation allows a reduction of the HAZ in the vertical and horizontal direction in comparison to contour welding. In order to quantify the influence of the beam and material properties on the HAZ geometries, a sensitivity analysis was performed. The simulation model is able to predict the HAZ height in good agreement to the experimental results. Here, the average deviation is 12.8%. However, there is a high deviation between the simulated and experimental HAZ with an average deviation of 48.5%. The reason for the high deviations will be investigated in future works. Possible reasons are simplifications, which were assumed in the simulation model. Another point is the influence of the interaction time between laser radiation and material on the melting temperature of the polymer. Past studies indicated that the melting temperature increases with decreasing interaction time.

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