A left and right truncated lognormal distribution for the stars

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Abstract The initial mass function for the stars is often modeled by a lognormal distribution. This paper is devoted to demonstrating the advantage of introducing a left and right truncated lognormal probability density function, which is characterized by four parameters. Its normalization constant, mean, the variance, second moment about the origin and distribution function are calculated. The chi-square test and the Kolmogorov–Smirnov test are performed on four samples of stars.

Keywords: Stars: characteristics and properties of Stars: normal

1 Introduction
The initial mass function (IMF) for the stars was firstly fitted with a power law by Salpeter, see [1]. He suggested \( p(m) \propto m^{-\alpha} \) where \( p(m) \) represents the probability of having a mass between \( m \) and \( m + dm \) and he found \( \alpha = 2.35 \) in the range \( 10M_\odot > M \geq 1M_\odot \). Secondly the IMF was fitted with three power laws, see [2,3,4] and four power laws, see [5,6,7]. The piecewise broken inverse power law IMF is

\[ p(m) \propto m^{-\alpha_i}, \quad (1) \]

each zone being characterized by a different exponent \( \alpha_i \) and two boundaries \( m_i \) and \( m_{i+1} \). In order to have a probability density function (PDF) normalized to unity, one must have

\[ \sum_{i=1,n} \int_{m_i}^{m_{i+1}} c_i m^{-\alpha_i} dm = 1 \quad . \quad (2) \]

The number of parameters to be found from the considered sample for the \( n \)-piecewise IMF is \( 2n - 1 \) when \( m_1 \) and \( m_{n+1} \) are the minimum and maximum of the masses of the sample. In the case of \( n = 4 \), which fits also the region of brown dwarfs (BD), see [8], the number of parameters is seven. In the field of statistical distributions, the PDF is usually defined by two parameters. Examples of two-parameter PDFs are: the beta, gamma, normal, and lognormal distributions, see [9]. The lognormal distribution is widely used in order to model the IMF for the stars, see [10,11,12,13]. The lognormal distribution is defined in the range of \( M \in (0, \infty) \) where \( M \) is the mass of the star. Nevertheless, the stars have minimum and maximum values, as an example from the MAIN SEQUENCE, an M8 star has \( M = 0.06M_\odot \) and an O3 star has \( M = 120M_\odot \), see [14]. The presence of boundaries for the stars makes attractive the analysis of a left and right truncated lognormal. In Section 2 the structure of the lognormal distribution is reviewed. In Section 3, the truncated lognormal distribution is derived. In Section 4.2 a comparison between the lognormal and truncated lognormal is done on four catalogs of stars. In Section 5 we compare the results of the truncated lognormal distribution with the double Pareto lognormal, the truncated beta, and the truncated gamma distributions.

2 The lognormal distribution
Let \( X \) be a random variable defined in \([0, \infty]\); the lognormal PDF, following [9] or formula (14.2)' in [15], is

\[ \text{PDF}(x; m, \sigma) = \frac{e^{-\frac{1}{2}(\ln(x) - m)^2}}{x \sigma \sqrt{2 \pi}}, \quad (3) \]
where $m$ is the median and $\sigma$ the shape parameter. The distribution function (DF) is

$$DF(x; m, \sigma) = \frac{1}{2} + \frac{1}{2} \text{erf} \left( \frac{1}{\sqrt{2}} \frac{\ln(x) - \ln(m)}{\sigma} \right),$$

where erf(x) is the error function, defined as

$$\text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt,$$

see [16]. The average value or mean, $E(X)$, is

$$E(X; m, \sigma) = me^{\frac{1}{2} \sigma^2},$$

the variance, $\text{Var}(X)$, is

$$\text{Var} = \sigma^2 \left( e^{\sigma^2} - 1 \right) m^2,$$

the second moment about the origin, $E^2(X)$, is

$$E(X^2; m, \sigma) = m^2 e^{2\sigma^2}.$$

The experimental sample consists of the data $x_i$ with $i$ varying between 1 and $n$; the sample mean, $\bar{x}$, is

$$\bar{x} = \frac{1}{n} \sum_{i=1}^{n} x_i,$$

the unbiased sample variance, $s^2$, is

$$s^2 = \frac{1}{n-1} \sum_{i=1}^{n} (x_i - \bar{x})^2,$$

and the sample $r$th moment about the origin, $\bar{x}_r$, is

$$\bar{x}_r = \frac{1}{n} \sum_{i=1}^{n} (x_i)^r.$$

The parameter estimation is here obtained in two ways. The matching moments estimator, (MME), is the first method:

$$E(X; m, \sigma) = \bar{x}_1; \quad E(X^2; m, \sigma) = \bar{x}_2,$$

and therefore

$$\hat{m} = \bar{x}_1 \frac{\sqrt{\bar{x}_2}}{\sqrt{\bar{x}_1}}$$

$$\hat{\sigma} = \sqrt{2} \sqrt{\ln \left( \frac{\sqrt{\bar{x}_2}}{\bar{x}_1} \right)}.$$

The second method implements the maximum-likelihood estimation (MLE), see [9].

3 The truncated lognormal distribution

Let $X$ be a random variable defined in $[x_l, x_u]$; the truncated lognormal PDF ($PDF_T$) is

$$PDF_T(x; m, \sigma, x_l, x_u) = \frac{\sqrt{2} e^{-\frac{1}{2} \frac{(\ln(x) - \ln(m))^2}{\sigma^2}}}{\sqrt{\pi} \sigma \left( \text{erf} \left( \frac{1}{2} \frac{\sqrt{2} \ln \left( \frac{x_u}{m} \right)}{\sigma} \right) - \text{erf} \left( \frac{1}{2} \frac{\sqrt{2} \ln \left( \frac{x_l}{m} \right)}{\sigma} \right) \right) x},$$

where $\text{erf}(x)$ is the error function, defined as

$$\text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt.$$
where \( m \) is now the scale parameter, \( \sigma \) is the shape parameter, \( x_l \) denotes the minimal value, and \( x_u \) denotes the maximal value. The introduction of the following coefficients allows a compact notation

\[
\begin{align*}
    a_1 &= \frac{1}{2} \sqrt{2} \left( -\sigma^2 + \ln(x_l) - \ln(m) \right), \\
    a_2 &= \frac{1}{2} \sqrt{2} \left( \sigma^2 + \ln(m) - \ln(x_u) \right), \\
    a_3 &= \frac{1}{2} \sqrt{2} \left( \ln(x_l) - \ln(m) \right), \\
    a_4 &= \frac{1}{2} \sqrt{2} \left( -\ln(x_u) + \ln(m) \right), \\
    a_5 &= \frac{1}{2} \sqrt{2} \left( -2\sigma^2 + \ln(x_l) - \ln(m) \right), \\
    a_6 &= \frac{1}{2} \sqrt{2} \left( 2\sigma^2 + \ln(m) - \ln(x_u) \right), \\
    a_7 &= \frac{1}{2} \sqrt{2} \left( -2\sigma^2 + \ln(x_u) - \ln(m) \right), \\
    a_8 &= \frac{1}{2} \sqrt{2} \left( \ln(x_u) - \ln(m) \right).
\end{align*}
\]

In the compact notation the PDF is

\[
PDF_T(x; m, \sigma, x_l, x_u) = \frac{-\sqrt{2}e^{-\frac{1}{2\sigma^2} \left( \ln(x) \right)^2}}{\sqrt{\pi} \sigma \left( \text{erf}(a_3) - \text{erf}(a_8) \right)} x,
\]

the DF is

\[
DF_T(x; m, \sigma, x_l, x_u) = \frac{-\text{erf} \left( \frac{1}{\sqrt{2}\sigma} \ln \left( \frac{x}{m} \right) \right) + \text{erf}(a_3)}{\text{erf}(a_3) - \text{erf}(a_8)},
\]

the mean, \( E(X)_T \), is

\[
E_T(X; m, \sigma, x_l, x_u) = \frac{e^{\frac{x}{\sigma}} \sigma^2 m \left( \text{erf}(a_1) + \text{erf}(a_2) \right)}{\text{erf}(a_3) + \text{erf}(a_4)},
\]

the variance, \( \text{Var}_T(X) \), is

\[
\text{Var}_T(X; m, \sigma, x_l, x_u) = \frac{N}{\left( \text{erf}(a_3) + \text{erf}(a_4) \right)^2},
\]

where

\[
N = e^{\sigma^2} \left( \text{erf}(a_3)\text{erf}(a_5)e^{\sigma^2} + \text{erf}(a_1)\text{erf}(a_6)e^{\sigma^2} + \text{erf}(a_4)\text{erf}(a_5)e^{\sigma^2} + \text{erf}(a_4)\text{erf}(a_6)e^{\sigma^2} \right) + \text{erf}(a_4)\text{erf}(a_6)e^{\sigma^2} - \left( \text{erf}(a_1) \right)^2 - 2\text{erf}(a_1)\text{erf}(a_2) - \left( \text{erf}(a_2) \right)^2 \right) m^2,
\]

the second moment about the origin, \( E_T^2(X) \), is

\[
E_T(X^2; m, \sigma, x_l, x_u) = -\frac{e^{2\sigma^2} m^2 \left( -\text{erf}(a_5) + \text{erf}(a_7) \right)}{\text{erf}(a_3) + \text{erf}(a_4)}.
\]

The two parameters \( x_l \) and \( x_u \) are the minimal and maximal elements of the sample. The two parameters \( m \) and \( \sigma \) can be found through the MME, first method

\[
E_T(X; m, \sigma, x_l, x_u) = \bar{x}_1 \quad ; \quad E_T(X^2; m, \sigma, x_l, x_u) = \bar{x}^2_2.
\]

The above system consists in two non-linear functions in two variables and can therefore be solved using the Powell hybrid method, see subroutine FORTRAN SNSQE in [17]. The second method implements the MLE in order to find \( m \) and \( \sigma \), see Appendix A.
4 Application to the stars

This section reviews some useful statistical parameters, such as the merit function $\chi^2$, the Akaike information criterion, and the Kolmogorov–Smirnov test. The four samples of stars which test the truncated lognormal distribution are introduced.

4.1 The adopted statistics

The merit function $\chi^2$ is computed according to the formula

$$\chi^2 = \sum_{i=1}^{n} \frac{(T_i - O_i)^2}{T_i},$$

where $n$ is the number of bins, $T_i$ is the theoretical value, and $O_i$ is the experimental value represented by the frequencies. The theoretical frequency distribution is given by

$$T_i = N \Delta x_i p(x),$$

where $N$ is the number of elements of the sample, $\Delta x_i$ is the magnitude of the size interval, and $p(x)$ is the PDF under examination. The size of the bins, $\Delta x_i$, is equal for each bin in the case of linear histograms, but different for each bin when logarithmic histograms are considered.

A reduced merit function $\chi^2_{red}$ is evaluated by

$$\chi^2_{red} = \chi^2 / NF,$$

where $NF = n - k$ is the number of degrees of freedom, $n$ is the number of bins, and $k$ is the number of parameters. The goodness of the fit can be expressed by the probability $Q$, see equation 15.2.12 in [18], which involves the degrees of freedom and $\chi^2$. According to [18] p. 658, the fit “may be acceptable” if $Q > 0.001$.

The Akaike information criterion (AIC), see [19], is defined by

$$AIC = 2k - 2ln(L),$$

where $L$ is the likelihood function and $k$ the number of free parameters in the model. We assume a Gaussian distribution for the errors and the likelihood function can be derived from the $\chi^2$ statistic

$L \propto \exp(-\frac{\chi^2}{2})$ where $\chi^2$ has been computed by eqn. (24), see [20], [21]. Now the AIC becomes

$$AIC = 2k + \chi^2.$$

The Kolmogorov–Smirnov test (K–S), see [22,23,24], does not require binning the data. The K–S test, as implemented by the FORTRAN subroutine KSONE in [18], finds the maximum distance, $D$, between the theoretical and the astronomical DF as well the significance level $P_{KS}$, see formulas 14.3.5 and 14.3.9 in [18]; if $P_{KS} \geq 0.1$, the goodness of the fit is believable.

4.2 The selected sample of stars

The test samples are selected from the Centre de Données astronomiques de Strasbourg (CDS) in order to ensure that the test can be easily reproduced, the name of the catalog is reported. The first test is performed on the low-mass IMF in the young cluster NGC 6611, see [25] and CDS catalog J/MNRAS/392/1034. This massive cluster has an age of 2–3 Myr and contains masses from $1.5 M_{\odot} > M \geq 0.02 M_{\odot}$. Therefore the brown dwarfs (BD) region, $\approx 0.2 M_{\odot}$ is covered. Table 1 shows the values of $\chi^2_{red}$, the AIC, the probability $Q$, of the fits and the two results of the K–S test: the maximum distance, $D$, between the theoretical and the astronomical DF as well the significance level $P_{KS}$. Figure 1 shows the fit with the truncated lognormal DF for NGC 6611, and Figure 2 the truncated lognormal PDF.
Figure 1. Empirical DF of mass distribution for NGC 6611 cluster data (207 stars + BDs) when the number of bins, $n$, is 20 (steps with full line) with a superposition of the truncated lognormal DF (dashed line). Theoretical parameters as in Table I MLE method. The horizontal axis has a logarithmic scale.

Figure 2. Frequencies of mass distribution for NGC 6611 cluster data (steps with full line) with a superposition of the truncated lognormal PDF (full line). Parameters as in Figure 1. The vertical and horizontal axes have a logarithmic scale.
Figure 3. Empirical DF of mass distribution for NGC 2362 cluster data (273 stars + BDs) when the number of bins, $n$, is 20 (steps with full line) with a superposition of the truncated lognormal DF (dashed line). Theoretical parameters as in Table 2, MLE method. The horizontal axis has a logarithmic scale.

Figure 4. Frequencies of mass distribution for NGC 2362 cluster data (steps with full line) with a superposition of the truncated lognormal PDF (full line). Parameters as in Figure 3. The vertical and horizontal axes have a logarithmic scale.
### Table 1. Statistical parameters of NGC 6611 (207 stars + BDs). The number of linear bins, $n$, is 20.

| PDF       | Method   | Parameters | $\alpha$ | $\sigma$ | $Q$ | $D$ | $AIC$ |
|-----------|----------|------------|----------|----------|-----|-----|-------|
| lognormal | MLE      | $\sigma=0.59$, $m=0.58$ | 11.84 | 2.96 | 0.09 | 0.0001 | 1.18 | 0.01 |
| lognormal | MLE      | $\sigma=0.59$, $m=0.58$ | 11.84 | 2.96 | 0.09 | 0.0001 | 1.18 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.59$, $m=0.58$, $x_{min}=54$ | 10.06 | 2.94 | 0.04 | 0.0005 | 2.39 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.59$, $m=0.58$, $x_{min}=54$ | 10.06 | 2.94 | 0.04 | 0.0005 | 2.39 | 0.01 |

The second test is performed on NGC 2362 where the 271 stars have a range $1.47M_\odot > M \geq 0.11M_\odot$, see [26] and CDS catalog J/MNRAS/384/675/table1. This is a very young open cluster with an estimated age of 3–9 Myr. Table 2 reports the statistical parameters, Figure 3 shows the fit with the truncated lognormal DF and PDF respectively.

### Table 2. Statistical parameters of NGC 2362 (272 stars). The number of linear bins, $n$, is 20.

| PDF       | Method   | Parameters | $\alpha$ | $\sigma$ | $Q$ | $D$ | $AIC$ |
|-----------|----------|------------|----------|----------|-----|-----|-------|
| lognormal | MLE      | $\sigma=0.69$, $m=0.68$ | 31.05 | 6.73 | 0.07 | 0.0004 | 4.02 | 0.01 |
| lognormal | MLE      | $\sigma=0.69$, $m=0.68$ | 31.05 | 6.73 | 0.07 | 0.0004 | 4.02 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.69$, $m=0.68$, $x_{min}=105$ | 33.04 | 6.71 | 0.07 | 0.0004 | 4.02 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.69$, $m=0.68$, $x_{min}=105$ | 33.04 | 6.71 | 0.07 | 0.0004 | 4.02 | 0.01 |

The third test is performed on a 40' circular field in the LMC made by 1563 stars in the range of masses, evaluated assuming an age of 4 Myr, $54M_\odot > M \geq 5M_\odot$, see [27] and CDS catalog J/ApJ/425/122/table2. Table 3 reports the statistical parameters. Figures 5 and 6 shows the fit with the truncated lognormal PDF.

### Table 3. Statistical parameters of a circular field in the LMC (1563 stars). The number of linear bins, $n$, is 20.

| PDF       | Method   | Parameters | $\alpha$ | $\sigma$ | $Q$ | $D$ | $AIC$ |
|-----------|----------|------------|----------|----------|-----|-----|-------|
| lognormal | MLE      | $\sigma=0.69$, $m=0.68$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |
| lognormal | MLE      | $\sigma=0.69$, $m=0.68$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.69$, $m=0.68$, $x_{min}=54$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.69$, $m=0.68$, $x_{min}=54$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |

The fourth test is performed on $\gamma$ Velorum cluster where the 237 stars have a range $1.31M_\odot > M \geq 0.15M_\odot$, see [28] and CDS catalog J/A+A/589/A70/table5. This cluster is consists of 5–10 Myr old premain sequence stars. The statistical parameters are reported in Table 4. Figures 7 and 8 shows the fit with the truncated lognormal DF and PDF respectively.

### Table 4. Statistical parameters of $\gamma$ Velorum cluster (237 stars). The number of linear bins, $n$, is 20.

| PDF       | Method   | Parameters | $\alpha$ | $\sigma$ | $Q$ | $D$ | $AIC$ |
|-----------|----------|------------|----------|----------|-----|-----|-------|
| lognormal | MLE      | $\sigma=0.69$, $m=0.68$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |
| lognormal | MLE      | $\sigma=0.69$, $m=0.68$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.69$, $m=0.68$, $x_{min}=54$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |
| truncated lognormal | MLE | $\sigma=0.69$, $m=0.68$, $x_{min}=54$ | 13.86 | 3.04 | 0.001 | 0.0001 | 1.08 | 0.01 |

5 Other new distributions

As an initial astronomical reference, we display a piecewise broken inverse power law PDF, see Figure 9.
Figure 5. Empirical DF of mass distribution for 1563 stars in LMC when the number of bins, \( n \), is 20 (steps with full line) with a superposition of the truncated lognormal distribution (dashed line). Theoretical parameters as in Table 3, MLE method. The horizontal axis has a logarithmic scale.

Figure 6. Frequencies of mass distribution for 1563 stars in LMC (steps with full line) with a superposition of the truncated lognormal PDF (full line). Parameters as in Figure 5. The vertical and horizontal axes have a logarithmic scale.
Figure 7. Empirical DF of mass distribution for 237 stars in \( \gamma \) Velorum cluster when the number of bins, \( n \), is 20 (steps with full line) with a superposition of the truncated lognormal DF (dashed line). Theoretical parameters as in Table 4, MLE method. The horizontal axis has a logarithmic scale.

Figure 8. Frequencies of mass distribution for 237 stars in \( \gamma \) Velorum cluster (steps with full line) with a superposition of the truncated lognormal PDF (full line). Parameters as in Figure 7. The vertical and horizontal axes have a logarithmic scale.
We now report three recent PDFs. The first is the double Pareto lognormal distribution which has PDF

\[ f(x; \alpha, \beta, \mu, \sigma) = \frac{1}{2} \alpha \beta \left( e^{\frac{1}{2} \alpha (\alpha \sigma^2 + 2 \mu - 2 \ln(x))} \text{erfc}\left( \frac{1}{2} \frac{\alpha \sigma^2 + \mu - \ln(x) \sqrt{2}}{\sigma} \right) \right) \]

\[ + e^{\frac{1}{2} \beta (\beta \sigma^2 - 2 \mu + 2 \ln(x))} \text{erfc}\left( \frac{1}{2} \frac{\beta \sigma^2 - \mu + \ln(x) \sqrt{2}}{\sigma} \right) x^{-1} (\alpha + \beta)^{-1}, \quad (29) \]

where \( \alpha \) and \( \beta \) are the Pareto coefficients for the upper and the lower tail, respectively, \( \mu \) and \( \sigma \) are the lognormal body parameters, and \( \text{erfc} \) is the complementary error function, see [29]. The mean (for \( \alpha > 1 \)) can be expressed as

\[ E(\alpha, \beta, \mu, \sigma) = \frac{\alpha \beta e^{\mu + \frac{1}{2} \sigma^2}}{(\alpha - 1)(\beta + 1)}. \quad (30) \]

This PDF exhibits a power law behaviour in both tails

\[ f(x) \sim k_1 x^{-\alpha-1} (x \to \infty) ; \quad f(x) \sim k_2 x^{\beta-1} (x \to 0), \quad (31) \]

where \( k_1 \) and \( k_2 \) are two constants. Figures 10 and 11 report the double Pareto lognormal DF and PDF respectively.

The second is the left truncated beta with scale PDF which is

\[ f_T(x; a, b, \alpha, \beta) = K x^{a-1} (b - x)^{\beta-1}, \quad (32) \]

where the constant is

\[ K = \frac{-a \Gamma (\alpha + \beta)}{b^{\beta-1} \Gamma a^\alpha \Gamma (\alpha + \beta) - b^{\beta-1+\alpha} \Gamma (1 + \alpha) \Gamma (\beta)}, \quad (33) \]

and

\[ H = _2F_1(\alpha, -\beta + 1; 1 + \alpha; \frac{a}{b}), \quad (34) \]

where \( _2F_1(a; b; c; z) \) is the regularized hypergeometric function [30], see [8]. Figure 12 reports the DF and Figure 13 the PDF.
Figure 10. Empirical DF of mass distribution for NGC 2362 cluster data (273 stars + BDs) when the number of bins, $n$, is 16 (steps with full line) with a superposition of the double Pareto lognormal DF (full line). Theoretical parameters as in Table 5. The horizontal axis has a logarithmic scale.

Figure 11. Frequencies of mass distribution for NGC 2362 cluster data (273 stars + BDs) when the number of bins, $n$, is 16 (steps with full line) with a superposition of the double Pareto lognormal PDF (full line). Theoretical parameters as in Table 5. The vertical and horizontal axes have a logarithmic scale.
Figure 12. Empirical DF of mass distribution for NGC 2362 cluster data (273 stars + BDs) when the number of bins, \( n \), is 20 (steps at full line) with a superposition of the left truncated beta DF (full line). Theoretical parameters as in Table 5. The horizontal axis has a logarithmic scale.

Figure 13. Frequencies for mass distribution in NGC 2362 cluster with (full line steps) and left truncated beta PDF. Parameters as in Figure 12. The vertical and horizontal axes have a logarithmic scale.
The third is the truncated gamma (TG) PDF which is

\[ f(x; b, c, x_l, x_u) = k \left( \frac{x}{b} \right)^{c-1} e^{-\frac{x}{b}} \]  

where the constant \( k \) is

\[ k = \frac{c}{b^c \Gamma(1 + c, \frac{x_u}{b}) - b^c \Gamma(1 + c, \frac{x_l}{b})} \]  

where

\[ \Gamma(a, z) = \int_z^\infty t^{a-1} e^{-t} dt, \]  

is the upper incomplete gamma function, see [31].

![Graph of DF vs Log10 (m/M⊙)](image)

**Figure 14.** Empirical DF of mass distribution for NGC 2362 cluster data (273 stars + BDs) when the number of bins, \( n \), is 18 (steps at full line) with a superposition of the truncated gamma DF (full line). Theoretical parameters as in Table 5. The horizontal axis has a logarithmic scale.

**Table 5.** Statistical parameters of NGC 2362 (272 stars) for different distributions

| PDF                | parameters | D  | \( P_{KS} \) |
|--------------------|------------|----|-------------|
| truncated lognormal | \( \sigma=0.59, \ m=0.625, \ x_l=0.119, \ x_u=1.47 \) | 0.047 | 0.556 |
| truncated gamma   | \( a=0.194, \ c=4, \ x_l=0.12, \ x_u=1.47 \) | 0.049 | 0.120 |
| double Pareto-lognormal | \( \alpha=5, \ \beta=2, \ \sigma=0.201, \ \mu_{LN}=-0.25 \) | 0.05 | 0.471 |
| left truncated beta | \( a=0.12, \ b=1.47, \ \alpha=2.18, \ \beta=2.93 \) | 0.048 | 0.53 |
| four inverse power law | \( m_1=0.11, \ m_2=0.22, m_3=0.41, m_4=0.78, m_5=1.47 \) | 0.084 | 0.052 |

**Figure 14** reports the truncated gamma DF and [Figure 15](#) the truncated gamma PDF. Table 5 reports the parameters of these three new PDFs as well as the parameters of the truncated lognormal in the case of NGC 2362. can be found Figure 16 displays all the PDFs here analysed.
Figure 15. Frequencies of mass distribution for NGC 2362 cluster data (273 stars + BDs) when the number of bins, $n$, is 18 (steps at full line) with a superposition of the truncated gamma PDF (full line). Theoretical parameters as in Table 5. The vertical and horizontal axes have a logarithmic scale.

Figure 16. Histogram (black step-diagram) of mass distribution as given by NGC 2362 cluster data (273 stars + BDs) with a superposition of the truncated gamma PDF (full red line), the lognormal PDF (dashed green line), the truncated lognormal PDF (dot-dash-dot-dash blue line), the truncated beta PDF (dotted cyan line) and the double Pareto lognormal PDF (dash-dot-dot-dot magenta line). Vertical and horizontal axes have logarithmic scales.
6 Conclusions

The truncated lognormal distribution gives better results, i.e. higher $P_{KS}$, than the lognormal distribution, see Tables 1, 2, 3 and 4 for the samples here considered. The lower and upper boundaries in mass are connected with the physical theories on the minimum and maximum mass for the stars. Fisher’s conjecture (see [32]) that statistical parameters are better inferred through the maximum likelihood estimator (MLE) than through the matching of moments estimator (MME) is also tested: in eight cases out of eight, the MLE produces better results, see Tables 1, 2, 3 and 4. The comparison of the truncated lognormal DF with other DFs assigns the best results to the truncated lognormal, i.e. higher $P_{KS}$, even if the difference from the double Pareto lognormal is small, see Table 5.

The number of free parameters of the truncated lognormal PDF is two once the lower and upper boundary are associated with the minimum and maximum mass of the considered sample, see Table 5 for the MLE method. In contrast, the number of parameters of the widely used four-piecewise broken inverse power law IMF is seven.

A The parameters of the truncated lognormal

The parameters of the truncated lognormal distribution can be obtained from empirical data by the maximum likelihood estimators (MLE) and by the evaluation of the minimum and maximum elements of the sample. Consider a sample $X = x_1, x_2, \ldots, x_n$ and let $x(1) \geq x(2) \geq \cdots \geq x(n)$ denote their order statistics, so that $x(1) = \max(x_1, x_2, \ldots, x_n)$, $(n) = \min(x_1, x_2, \ldots, x_n)$. The first two parameters $x_l$ and $x_u$ are

$$x_l = x(n), \quad x_u = x(1).$$

(38)

The MLE is obtained by maximizing

$$A = \sum_i^n \ln(TL(x_i; m, \sigma, x_l, x_u)).$$

(39)

The two derivatives $\frac{\partial A}{\partial m} = 0$ and $\frac{\partial A}{\partial \sigma} = 0$ generate two non-linear equations in $m$ and $\sigma$ which can be solved numerically, we used FORTRAN subroutine SNSQE in [17].

$$\frac{\partial A}{\partial m} = (\text{erf}(\frac{1}{2} \frac{\sqrt{2}(\ln(x_l) - \ln(m))}{\sigma}) - \text{erf}(\frac{1}{2} \frac{\sqrt{2}(\ln(x_u) - \ln(m))}{\sigma}))$$

$$-\text{erf}(\frac{1}{2} \frac{\sqrt{2}(\ln(x_u) - \ln(m))}{\sigma})(\ln(m) - \sum_{i=1}^n \ln(x_i))) = 0 ,$$

(40)

and

$$\frac{\partial A}{\partial \sigma} = \frac{N}{D} = 0,$$

(41)
where

\[
N = \ln(x_u) \sqrt{2e}^{-\frac{1}{2} \frac{\ln(x_u) - \ln(m))^2}{\sigma^2}} n\sigma - \ln(x_l) \sqrt{2e}^{-\frac{1}{2} \frac{\ln(x_l) - \ln(m))^2}{\sigma^2}} n\sigma
\]

\[
+ \sqrt{2e}^{-\frac{1}{2} \frac{\ln(x_u) - \ln(m))^2}{\sigma^2}} \ln(m)n\sigma - \sqrt{2e}^{-\frac{1}{2} \frac{\ln(x_l) - \ln(m))^2}{\sigma^2}} \ln(m)n\sigma
\]

\[
+ n(\ln(m))^2 \sqrt{\pi} \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_u) - \ln(m))}{\sigma}\right)
\]

\[
- n\sigma^2 \sqrt{\pi} \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_l) - \ln(m))}{\sigma}\right)
\]

\[
- n(\ln(m))^2 \sqrt{\pi} \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_l) - \ln(m))}{\sigma}\right)
\]

\[
+ n\sigma^2 \sqrt{\pi} \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_l) - \ln(m))}{\sigma}\right)
\]

\[
\sum_{i=1}^{n} \ln(x_i)(\ln(x_i) - 2 \ln(m)) \sqrt{\pi} \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_i) - \ln(m))}{\sigma}\right)
\]  

\[
- \sum_{i=1}^{n} \ln(x_i)(\ln(x_i) - 2 \ln(m)) \sqrt{\pi} \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_l) - \ln(m))}{\sigma}\right)
\],

(42)

\[
D = \sqrt{\pi} \left( - \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_l) - \ln(m))}{\sigma}\right) \right)
\]

\[
+ \text{erf}\left(\frac{1}{2} \frac{\sqrt{2}(\ln(x_u) - \ln(m))}{\sigma}\right) \right) \sigma^3
\].

(43)

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