Inkjet-Printed Micrometer-Thick Perovskite Solar Cells with Large Columnar Grains

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Transferring the high power conversion efficiencies (PCEs) of spin-coated perovskite solar cells (PSCs) on the laboratory scale to large-area photovoltaic modules requires a significant advance in scalable fabrication methods. Digital inkjet printing promises scalable, material, and cost-efficient deposition of perovskite thin films on a wide range of substrates and in arbitrary shapes. In this work, high-quality inkjet-printed triple-cation (methylammonium, formamidinium, and cesium) perovskite layers with exceptional thicknesses of >1 μm are demonstrated, enabling unprecedentedly high PCEs >21% and stabilized power output efficiencies >18% for inkjet-printed PSCs. In-depth characterization shows that the thick inkjet-printed perovskite thin films deposited using the process developed herein exhibit a columnar crystal structure, free of horizontal grain boundaries, which extend over the entire thickness. A thin film thickness of around 1.5 μm is determined as optimal for PSC for this process. Up to this layer thickness X-ray photoemission spectroscopy analysis confirms the expected stoichiometric perovskite composition at the surface and shows strong deviations and inhomogeneities for thicker films. The micrometer-thick perovskite thin films exhibit remarkably long charge carrier lifetimes, highlighting their excellent optoelectronic characteristics. They are particularly promising for next-generation inkjet-printed perovskite solar cells, photodetectors, and X-ray detectors.

1. Introduction

Within the last decade, hybrid organic–inorganic lead halide perovskite semiconductors have attracted enormous attention in science and technology for lighting-emitting devices,[1–3] lasers,[4,5] photodetectors,[6] X-ray detectors,[7–9] and, in particular, photovoltaics (PVs).[10–14] Recently, power conversion efficiencies (PCEs) exceeding 25% have been demonstrated in perovskite thin film solar cells (PSCs).[15] The rapid advance of the multocrystalline perovskite thin films builds on excellent optoelectronic properties, namely, strong absorption coefficients,[16] excellent tolerance to defects, and high charge carrier mobilities.[17–20] Moreover, given the exceptional ability to tune the band gap of perovskite semiconductor by compositional engineering,[4,21–23] this class of materials is perfectly suited for realizing tandem solar cells.[24]

Solution processing of multocrystalline thin films offers an easy and energy-efficient route to manufacture perovskite photovoltaics, given the low crystallization temperature of the perovskite thin films compared to other photovoltaic technologies. Up to date, the research and development of PSCs has been mostly focused on laboratory scale, particularly on spin-coated perovskite thin films.[10–12,14]

However, with view to a future commercialization of the technology, scalable fabrication of perovskite PV remains a key challenge.[25] In recent years, research and development of solution-based scalable coating and printing techniques have attracted significant attention, for example, blade coating.[26,27]
slot-die coating,[28,29] and inkjet printing.[30–35] Inkjet printing stands out among these technologies as a noncontact digital printing technique that offers the freedom of printing arbitrary design patterns[37] at very low material consumption.[38] Inkjet printing is an adaptable and fast printing technique and is used not only in research on, for example, large-area organic solar cells,[39] but also already in large-area industrial applications such as next-generation organic light-emitting diode production lines.[40,41] Over the past years, (partly) inkjet-printed (IJP) PSCs such as next-generation organic light-emitting diode production procedures.[30,32,33,44–46] In order to close the gap in performance, three key challenges need to be addressed regarding the processing of IJP PSCs: First, achieving optimal printability of perovskite thin films by inkjet printing requires engineering of the solvent composition, whereby all the precursor materials are dissolved in a so-called ink.[47–49] Second, the underlying charge transport layer not only needs to enable optimal charge carrier extraction, but it also needs to support optimal wetting of the printed ink droplets. This optimization requires ongoing engineering of the ink–surface interaction (e.g., by means of surface treatment) along with optimizations of the layer stack.[36,46] Third, for high performance PSCs, the nucleation and crystallization of the deposited perovskite wet films needs to be controlled.[50,51] For spin-coated perovskite thin films, the so-called anti-solvent treatment is an established strategy to control crystallization of multi-cation perovskite thin films.[14,52] However, given the large amounts of solvents required and the very complex timing of the prompt crystallization initiated by the anti-solvent treatment, it needs to be modified or replaced to realize scalable deposition, for example, by a vacuum-assisted crystallization.[33,53]

In this work, these three challenges are addressed with the objective to optimize PCE and stability of IJP PSCs. In particular, the ink composition, surface treatment, deposition, and annealing methods of IJP perovskite films are presented for different printing parameters. Material characterization reveals the crystal growth, lead iodide content, surface composition, and charge carrier lifetimes. IJP multicrystalline perovskite absorber layers with large columnar crystals are deposited, exhibiting grains extending over the entire layer thickness. With increasing printing resolution, hence thickness, the perovskite films demonstrate lower relative lead iodide content and longer charge carrier lifetimes. Interestingly, PSCs fabricated using IJP perovskite thin film absorbers (example stack in Figure 1a) perform better for layer thicknesses significantly higher than for comparable state-of-the-art spin-coated devices.

Champion devices with a thickness of around 1.5 μm exhibit a single scan PCE of >21% (Figure 1b), a stabilized PCE of up to 18.5% (Figure 1c), and show no significant drop in PCE after 72 h (Figure 1e). Throughout the detailed characterization of the underlying nickel oxide (NiOx) layer the stabilized PCE has been published earlier this year.[46] The demonstrated short-circuit current-density (Jsc) of these IJP PSCs is in the range of the highest Jsc reported for spin-coated PSCs of similar architecture and same band gap[10,12,13] as well as, to the authors’ best knowledge, considerably higher than any IJP PSCs.[25] This improvement originates from exceptionally thick absorber layers, compared to state-of-the-art spin-coated and excellent charge carrier lifetimes. In addition, open-circuit voltage (Voc) and fill factor are remarkably high, inter alia indicating that the IJP perovskite thin film exhibits high diffusion lengths. The PSCs demonstrated in this work show low hysteresis (champion device’s hysteresis index factor (HIF) is 0.84; see the Experimental Section), which is remarkable for IJP PSCs.[31,46] The presented results are the highest stabilized PCE for IJP PSCs and both PCE and stabilized PCE are considerably closer to state-of-the-art values achieved with spin-coating than any preceding reports (Figure 1d).[25,32,54,55]

2. Manufacturing of Inkjet-Printed Solar Cells

2.1. Fabrication of Perovskite Solar Cells

The perovskite solar cells reported in this study are processed in the p-i-n architecture based on the layer sequence indium tin oxide (ITO)/NiOx/perovskite/C60/bathocuproine (BCP)/gold as depicted in Figure 1a (for a detailed description of the fabrication process see the Experimental Section). Prepatterned ITO-coated glass substrates were coated with NiOx as hole transport layer (HTL) deposited by electron beam evaporation as described in detail by Abzieher et al.[36] Subsequently, the triple-cation mixed halide perovskite (TCP) absorber layers were deposited. The reference spin-coated samples were prepared using the widely established anti-solvent step and consecutive annealing on a hotplate in an inert nitrogen atmosphere.[52,56] The IJP samples were processed under controlled ambient conditions (relative humidity of ~45%). The key steps of the perovskite thin film are depicted in Figure 2 and described in detail in the following section. The solar cells were completed by evaporating a C60 fullerene electron transport layer on top, followed by a 3 nm thin BCP interfacial layer. Finally, a gold rear electrode is thermally evaporated using a shadow mask, which defines the active area to 10.5 mm².

2.2. Inkjet-Printed Triple-Cation Perovskite Thin Films

Fabrication of high performance IJP perovskite thin films involves three key steps, as shown in Figure 2: (1) Ink preparation: the ink needs to be engineered such that ink droplets of well-defined size and shape are generated with a given ink system; (2) Inkjet printing: the droplets have to be printed on the substrate, where the forces between ink (cohesion of the droplet) and substrate (adhesion to the surface) need to be balanced to avoid repelling on the one side and uncontrolled
spreading over the whole substrate on the other side; and (3) Drying and annealing: the solvents in the as-printed wet films need to evaporate and the remaining precursor materials need to crystallize in a pinhole-free perovskite thin film.

2.2.1. Ink Preparation

As the printing setup used in this work only allows for a single channel print process, a single ink approach is presented. Although without multichannel the perovskite composition cannot be changed during printing, the single channel approach makes the process accessible to a large variety of inkjet printing setups with similar constraints. Hence, we prepare a single ink for deposition of TCP thin films, containing cesium (Cs), methylammonium (CH$_3$NH$_3$; MA), and formamidinium (CH(NH$_2$)$_2$; FA), similar to Mathies et al.\cite{46} The targeted stoichiometric ratio of the TCP perovskite thin films is Cs$_{0.1}$MA$_{0.15}$FA$_{0.75}$Pb$_{0.85}$Br$_{0.15}$. For this, the precursor materials (see the Experimental Section) are dissolved in a mixed solvent system. In our studies a mixture of the high-boiling-point solvent γ-butyrolactone (GBL) and the polar-aprotic solvents N,N-dimethylformamide (DMF) as well as dimethyl sulfoxide (DMSO) resulted in improved homogenous drying.

2.2.2. Inkjet Printing

The droplets ejected from the print-head of the inkjet printing system impinge next on the substrate. In inkjet printing wetting behavior of the droplets on the substrate surface is classified as de-wetting, over-wetting, and optimal wetting. While the first describes the contraction of the as-printed droplets on the surface without formation of a continuous wet film, the second describes the inhomogeneous spread of droplets across the substrate with in the worst case no pinning behavior. In order to achieve optimal (partial) wetting behavior, a minimum contact angle $\theta$ of about 5°–10° between the droplet and the substrate is desired.
substrate surface is needed to avoid over-wetting. On the other hand, \( \theta < 90^\circ \) is the theoretical upper limit, since the droplets start to repel from the surface, which leads to de-wetting. To tune the contact angle, the surface tension (SFT) of the ink and the surface free energy (SFE) of the substrate need to be adjusted. To estimate the wettability of the TCP ink on the NiO\textsubscript{x} HTL used in this work, the SFT shifts from the wetting envelope toward higher values, the larger the contact angle, leading ultimately to an entire de-wetting. In agreement with the depicted wetting envelopes in Figure 3, the NiO\textsubscript{x} HTL used in this work leads to contact angles \( \theta > 25^\circ \) that imply small droplet diameters.
Although it is possible to print on surfaces with contact angles of $\theta \approx 25^\circ$, often an increased density of (pin-)holes due to local de-wetting (see Figure S2 in the Supporting Information) is observed for this range of angles for IJP perovskite thin films. To minimize the de-wetting, a short low-power oxygen plasma is applied prior to the printing process to increase the overall SFE in a way that the SFT of the ink lies inside the envelope. Although this does not exactly match the theory (it should be $\theta_{\text{theory}} < \theta_{\text{envelope}}$ but it is $\theta \approx 10^\circ > \theta_{\text{envelope}}$), the wetting improves, which results in larger drop diameters and increases the feature size ($\approx 100 \mu m$ instead of $\approx 70 \mu m$; see Figure 3b), and thus reduces the risk of formation of holes in the wet film significantly (see Figure S2 in the Supporting Information). Although the demonstrated droplet diameters of 100 $\mu m$ are rather large, we are still able to print almost arbitrary patterns even on relatively small scales with sub-millimeter structures with this process (Figure S3 in the Supporting Information).

2.2.3. Drying and Annealing

Subsequent to the printing of the wet film, the samples are moved manually to a nearby vacuum chamber. Since the wet film is of significant thickness ($\approx 20 \mu m$ for 1100 dots per inch (dpi), calculated), the samples have to be moved with care to avoid movement of the wet-film profile. The vacuum chamber is evacuated for a few minutes. While the pressure is going down to $5 \times 10^{-2}$ mbar, the boiling point of the solvents decreases and the evaporation rate of the solvents increases. According to the ascending boiling point, DMF starts to evaporate first, followed by DMSO and GBL. While the solvents are getting extracted, the shrinking wet film starts to crystallize, indicated by a change in color from yellow to dark brown. After that, the chamber is slowly vented with ambient air and the samples are subsequently annealed on a hotplate under ambient air.

3. Results

3.1. Inkjet-Printed Perovskite Thin Films with Large Columnar Grains

The PSCs with IJP multicrystalline perovskite absorber layers introduced in Figure 1 exhibit an exceptional thickness around 1.5 $\mu m$ along with high PCEs and stable power outputs. In order to investigate the morphology and the optoelectronic characteristics of the micrometer-thick IJP perovskite layers, their thickness along with the printing resolution—given in dpi—is varied.

Scanning electron microscopy (SEM) images of the cross-section of the perovskite layers, shown in Figure 4a, provide an estimate of the respective thicknesses, which are in good agreement with the profilometer measurements that are used to determine the actual thickness. The layer thickness increases approximately quadratic with the printing resolution, as expected by theory and described in previous reports.\cite{46} The perovskite thin film thickness ranges from $\approx 400 \text{ nm}$ (600 dpi) up to almost 4 $\mu m$ (2000 dpi). It should be noted that such
thick perovskite thin films cannot be deposited by spin-coating, where spin speed and solubility limit the thickness. Much thinner layers are also commonly used for state-of-the-art PSCs processed with alternative scalable fabrication techniques, e.g., 500–600 nm for blade coating,[26,60] 350–430 nm for slot-die coating,[29,61] or 350 nm for thermal co-evaporation.[36,62] In previous reports on PSCs with perovskite absorber thin films deposited by inkjet printing, horizontal grain boundaries are apparent, introducing defects that induce nonradiative recombination and impede the charge transport through the absorber.[46] In contrast, the IJP perovskite thin films of this work exhibit large columnar grains extending over the entire perovskite thin film even for the thickest layers (over 4 µm) examined in this work (see SEM images of the layer cross-sections in Figure 4a). This allows for perovskite absorber layers with only few to no horizontal grain boundaries, improving its transport properties. SEM images of the perovskite surface indicate an increasing grain size for higher resolutions in horizontal directions (compare Figure S5 in the Supporting Information), suggesting a similar trend as observed for the formation of the vertical grains (Figure 4a).

To further investigate the influence of inkjet printing with different printing resolution, as well as the differences in annealing on the formation of the perovskite thin films, stacks of glass/ITO/NiO$_x$ and perovskite were investigated by X-ray diffraction (XRD) (Figure 4b). The XRD diffractograms exhibit diffraction peaks for the perovskite in both the printed layers and the spin-coated reference, including the prominent diffraction peak of the (001)-plane of the cubic perovskite phase at 14.2° and the combined peak of the (012) and (102)-plane at ≈32°.[63] However, the relative intensity of the peaks assigned to different crystal planes of the perovskite crystal structure systematically changes with the thickness of the perovskite thin film. For example, compared to the (001)-peak, the (012)-(102)-peak grows disproportionately from a 4:1 ratio for 600 dpi to less than a 1:1 ratio for 2000 dpi (2:5:1 for the spin-coated reference). This indicates that crystallization along either the (012) or (102)-plane is increasingly dominant with thicker layers supporting the observation of a change in preferred crystal orientation made by SEM.

Not only the preferred crystal orientation, but also the overall composition of the perovskite changes: spin-coated layers show a significantly lower lead-iodide-to-perovskite ratio than printed layers of comparable thickness (600–800 dpi), which can be determined by comparison of the diffraction peak attributed to lead iodide (12.8°) and the diffraction peak attributed to the perovskite (100)-plane (see Figure 4b). This indicates an enhanced lead iodide volume fraction in the IJP perovskite layers, which is attributed to the differences in the fabrication process, for example a significantly slower vacuum annealing step compared to anti-solvent treatment for spin-coated perovskite thin films as well as a modified solvent system. Furthermore, it should be noted that the lead-iodide-to-perovskite ratio decreases with increasing printing resolution. On the one hand, this could indicate an overall different crystal composition, for example, due to slower drying of the thicker wet films printed with high resolutions. On the other hand, this could also point to local differences in crystallization: since higher printing resolution entail thicker thin films this is an indication of lead-iodide-rich layers forming at a perovskite interface (with either the underlying HTL or air) that decreases in volume fraction with increasing thin film thickness.

It should be further noted that the roughness of IJP perovskite thin films increases with printing resolution and is significantly higher compared to spin-coated thin films, due to the absence of a centrifugal force and a slower drying process (see above) giving rise to fluid mechanically driven artifacts, such as coffee rings (compare Figure S4 in the Supporting Information). In order to investigate the composition at the surface of these IJP perovskite thin films, spatially resolved X-ray photoemission spectroscopy (XPS) mapping was conducted. The investigated layer stack represents an uncompleted PSC (glass/ITO/NiO$_x$/TCP). The XPS maps display the atomic ratios of cesium, nitrogen (originating from the MA and FA cations), iodine, and bromine in comparison to lead over an area of 8 x 8 mm$^2$ (Figure 5a). The more uniformly colored maps observed for printing resolutions of 600–1100 dpi indicate fewer local differences in the surface composition of the films and therefore, an overall more homogeneous layer. Figure 5b summarizes the distribution of the element ratios (Cs/Pb, N/Pb, Br/Pb, and I/Pb) measured at the perovskite surface. The decrease in homogeneity of the perovskite composition at the surface with increasing printing resolution is apparent as a broadening of all elemental ratio distributions. Not only are the perovskite thin-film surfaces printed with 600–1100 dpi the most homogeneous, but for these layers the average composition of the surface also deviates less from the desired perovskite composition of Cs$_6$Pb$_{2}$Br$_{17}$,FA$_{1.75}$Pb$_{1.5}$I$_{2}$Br$_{1.15}$. In this context it should be noted that for all element ratios a small lead deficiency is observed, which is in agreement to other XPS studies.[64–68] Very thick (＞1.5 µm) perovskite thin films processed with higher resolution (1400 and 2000 dpi) are significantly more inhomogeneous and contain significant excess of halides, cesium, and organic cations at the surface than the films deposited at lower resolution.

Having investigated the stoichiometric composition and crystallinity, next the good optoelectronic quality of the thick IJP perovskite thin films is highlighted by photoluminescence (PL) measurements. A two-term exponential fit (see the Experimental Section for details) was used to determine a mixed time constant τ$_1$ and a charge carrier lifetime constant τ$_2$ as decay constants of time-resolved PL in perovskite thin films on glass/ITO/NiO$_x$ (Figure 6).[69–72] While the charge extraction time constant is similar for all samples, layers printed with higher resolutions (especially 1100 dpi and more) display significantly longer charge carrier lifetime constants, emphasizing the excellent optoelectronic quality. This is attributed to the formation of large columnar grains extending over the entire thin film thickness, minimizing the influence of crystal domain grain boundaries that are prone to recombination centers and non-radiative recombination. The comparison of spin-coated to low resolution-printed films with comparable thickness (600 dpi) shows similar lifetime constants, which suggests that the changes in relative lead iodide content seen in XRD are in this case not detrimental for the changed charge carrier transport.[71] A comparison with PL measurements on perovskite layers inkjet printed on glass (Figure S6, Supporting Information) suggests that while quenching at the NiO$_x$ interface is an
observable effect, it is not the dominant mechanism behind the change of lifetime constant $\tau_2$. It should be highlighted that the charge carrier lifetime constant in the thicker perovskite thin films (1100–2000 dpi) is in the range of 0.5–1 µs, which is a truly remarkable value in view of lifetimes reported previously for multicrystalline perovskite thin films (up to values of 480 ns for triple-cation perovskites, including determination with a two-term exponential fit\[72,73\]) and proving the exceptional optoelectronic quality of printed perovskite films presented in this work.

**Figure 5.** a) Local distribution of element ratios at the perovskite surface determined by spatially resolved X-ray photoemission spectroscopy (XPS) mapping of IJP perovskite on glass/ITO/NiO$_x$ stacks. Mapping area is 8 × 8 mm$^2$ for each sample. Different printing resolutions for the IJP perovskite are displayed in the columns. Four elemental ratios (Cs/Pb, N/Pb, Br/Pb, and I/Pb) are shown in the rows. b) Statistical analysis of the XPS maps.

### 3.2. Perovskite Solar Cells with Inkjet-Printed Layers

Considering the long lifetimes of the thick IJP perovskite layers with large columnar grains, the performance of IJP PSCs is investigated next. Although theoretical expectations of diffusion lengths for electrons and holes in perovskite range up to 1 µm and above,\[20\] for most state-of-the-art spin-coated PSCs an absorber thickness of around 400–600 nm has proven to result in most efficient devices. Likewise, many works published so far on PSCs via inkjet printing report...
PCE-maxima for similar thickness. Therefore, the performance of IJP PSC with different layer thicknesses, achieved by variation of printing resolution, is compared here. With increasing printing resolution, thus thickness, an optimal point of performance becomes apparent at 1100 dpi (roughly 1.5 µm layer thickness), which can be related to the material and layer properties investigated by SEM, PL, XRD, and XPS measurements (Figure 7): While the fill factor is the highest for devices printed with low resolution, the short-circuit current density $J_{sc}$ increases with resolution up to a maximum of 24 mA cm$^{-2}$ for 1100 dpi, respectively 1.5 µm, which is similar to a calculated thickness with maximum $J_{sc}$ for low recombination regime MAPbI$_3$ of 1.3 µm. As detailed above, the increase in $J_{sc}$ in the devices is attributed to the excellent charge carrier lifetimes measured and the low number of horizontal grain boundaries of the columnar crystal grains extending over the thin film (compare Figure 4a) and increases significantly up to this resolution. Additionally, with increasing resolution the volume ratio of lead iodide compared to perovskite in the thin film is decreasing, which is assumed to be beneficial for the device. For layers printed with higher resolutions (1400–2000 dpi), both current density and fill factor decrease significantly (Table S7, Supporting Information). According to the material characterization two possible reasons for this can be identified: First, such high printing resolutions entail perovskite thin film thicknesses of up to 4 µm. In spite of the exceptional long charge carrier lifetimes, efficient charge transport over the entire thickness might not be possible anymore. Second, the perovskite thin-film surface gets less homogeneous and differentiates more
from the desired stoichiometric composition influencing the charge carrier extraction.

There is an optimal printing resolution for IJP PSC that results from a trade-off between: i) increasing lifetime as well as decreasing lead-iodide ratio; and ii) increasing thickness and decreasing surface quality. The maximum PCE is identified for devices with IJP perovskite layers with a printing resolution of about 1100 dpi. This corresponds to a thickness of roughly 1.5 µm. Simple drop-coating experiments with similar ink-volume, which should mimic the deposition by inkjet printing, demonstrate a similar trend, affirming a correlation between PCE and thin film thickness. However, a wide statistical spread in performance of drop-coated devices compared to IJP ones is observable, presumable since especially area and therefore wet-film thickness is harder to control. Best IJP devices allow for maximum PCEs of over 21% in backward J–V characteristics and average PCEs of almost 20% (best: 19.9%). To the best of the authors’ knowledge, this is the highest PCE for IJP perovskite layers so far. The hysteresis effect is, on average, lowest for devices fabricated with 1100 dpi printing resolution (compare HIF in Table S7 in the Supporting Information). However, it is not negligible. Hence, an average PCE or as chosen here the stabilized PCE under maximum power point (MPP) tracking is not negligible. Hence, an average PCE or as chosen here the stabilized PCE under maximum power point (MPP) tracking is more meaningful to evaluate the solar cell performance. The maximum short-time stabilized PCE is 18.5%. High-performance PSCs were reproducibly fabricated over seven batches with an average single-scan PCE of 16.3% ± 2.5% for 178 solar cells (see Figure S9 in the Supporting Information). Although the average is considerably lower than the best PCE, this is a result with significant impact, since addressing the problems with reproducibility and stability of devices is essential for commercialization not only for IJP PSCs but also for PSCs in general. Here, solar cells with thick absorber layers might also be of advantage, since the probability of short-circuits through formation of pinholes significantly decreases with thickness leading to fewer shunted devices per batch and assumable to reduced material degradation facilitated by pinholes. To provide proof-of-concept for scalability, an immature prototype of a PSC with active area of ≈1 cm² was fabricated (see Figure S10 in the Supporting Information). The mean PCE of 10.4% is mainly limited by a low fill factor, which is assumed to origin in a non-optimized layout. $V_{OC}$ and $J_{SC}$ are comparable to the standard devices.

### 3.3. Conclusion

Here, we demonstrate IJP PSCs with a PCE of over 21% and also exhibiting >18% stabilized PCE—both of these values are the highest reported for partly IJP PSCs so far—as well as good reproducibility. Champion devices incorporate very thick absorber layers of about 1–1.5 µm (printing resolution 1100 dpi). The IJP perovskite thin films exhibit columnar perovskite crystal structures that extend over the entire thickness without horizontal grain boundaries, as evidenced via scanning electron microscopy. These columnar perovskite crystal structures are assumed to have few defects and facilitate efficient charge carrier transport. In addition, the printed perovskite thin films show long charge carrier lifetimes of >0.5 µs, proven by time-resolved photoluminescence measurements. Although an effective decrease of relative PbI$_2$ amount in the perovskite for thicker layers by X-ray diffraction is observable, X-ray photoemission spectroscopy studies indicate that only layers up to ≈1.5 µm thickness exhibit the desired surface composition with a high degree of homogeneity. This work demonstrates that the PSCs with IJP triple-cation absorber can be close to PCE of state-of-the-art devices by other deposition methods, enabling this material-efficient process, which is well suited for upscaling while additionally offering the possibility to process in almost arbitrary patterns.

In view of the future development of the technology, this study provides a remarkable demonstration that very thick perovskite absorber layers that are possibly easier to scale up without defects can be deposited at high quality by inkjet printing.

### 4. Experimental Section

**Sample Fabrication:** For preparation of the samples, the following route was used: Prepatterned ITO on glass substrates (Luminescence Technology) were cleaned consecutively in acetone and isopropanol in an ultrasonic bath for 10 min each, followed by an oxygen plasma cleaning step for 2–3 min. Afterward 20 nm NiO$_x$ HTL was deposited as described in detail by Abzieher et al.\[46] Basically, green raw material NiO$_x$ lumps (Alfa Aesar, 99.995%) were filled into an alumina crucible in an electron evaporation system (Angstrom Engineering) and were evaporated at base pressures around $10^{-6}$ mbar, while using a small (about 10 sccm) oxygen flow to maintain the original composition ratio.

Subsequently, the TCP absorber layers were deposited. The ink for inkjet printing was prepared as follows: CH(NH$_2$)$_2$I (FAI, 0.6 µm, GreatCell Solar), PbI$_2$ (0.66 µm, TCI Chemicals), CH$_3$NH$_3$Br (MABr, 0.12 µm, GreatCell Solar), and PbBr$_2$ (0.12 µm, TCI Chemicals) were dissolved in a mixture of DMF (Sigma-Aldrich), DMSO (Merck), and GBL (Sigma-Aldrich) in a ratio of 28:26:46 (volume percentage). Additionally, CsI (1.5 µm, Alfa Aesar) was dissolved in DMSO and then added to the first solution to get a 0.75 µm TCP solution with the composition Cs$_{0.10}$FA$_{0.75}$MA$_{0.15}$Pb(Br$_{0.15}$I$_{0.85}$)$_2$. Before printing, the ink was filtered with a 0.45 µm pore size polytetrafluoroethylene (PTFE) filter. The triple-cation solution for spin coating was prepared with the same precursor salts (FAI 1 µm, PbI$_2$ 1.1 µm, MABr 0.2 µm, and PbBr$_2$ 0.2 µm) in a mixture of DMF and DMSO in a ratio of 4:1 (volume percentage). The same CsI 1.5 µm solution was then added to achieve a C$_{58.10}$FA$_{73}$MA$_{11}$Pb(Br$_{0.15}$I$_{0.85}$)$_2$ composition.

For inkjet printing of the TCP layers, a Meyer Burger PixDRO LP50 with a print head module for 10 pl Fujifilm cartridges (Dimatix DMC-16610) was used. Prior to the printing, the NiO$_x$ surface was treated with a short oxygen plasma at low power. For every printed resolution, a jetting frequency of 2 kHz was used together with a single pulse waveform with a peak voltage of 33 V and a pulse width of 5 µs. The total area of inkjet-printed TCP was 12 × 12 mm² per sample. Within the time frame of 30 s the as-printed samples were moved by hand to a nearby vacuum chamber (Pfeiffer Vacuum Technology AG) which was evacuated down to about 5 × 10⁻² mbar. Then, the chamber was slowly vented and the samples were annealed on a hotplate at 100 °C. The complete printing procedure was done in ambient atmosphere (≈23 °C, ≈45% relative humidity). For spin coating of the TCP layers, the NiO$_x$ samples were also treated with a short oxygen plasma and then moved into a nitrogen filled glove box. The layer was spin-coated using a two-step program (1000 rpm for 10 s, 5000 rpm for 20 s) with a chlorobenzene (Sigma-Aldrich) anti-solvent step (100 µL) and then annealed for 1 h on a hotplate at 100 °C.

For completing the p-i-n-solar cell stack, a 25 nm thick C$_6$O$_{12}$ fullerene electron layer, followed by a 3 nm thick BCP interfacial layer, was thermally evaporated on top of the perovskite layer. Finally, a 60 nm
thick gold back contact was thermally evaporated using a shadow mask, which defined the active area to 10.5 mm² per solar cell with four cells per sample.

Characterization—SEM: SEM images were taken with a Zeiss Auriga system. Measurements were taken in high vacuum. For cross-sections, samples were cut on the rear side and broken. Pictures were not corrected for small tilts or errors due to breaking. Since the integration time of pictures may also have a small effect (due to shifts caused by electrical charging and discharging) on the displayed scale, these pictures were not reliable for thickness determination.

Solar Cell Characterization: For the measurement of the solar cell characteristics, a xenon-lamp-based solar simulator (Newport Oriel Sol3A) inside a nitrogen filled glove box with an AM1.5 spectrum (100 mW cm⁻²) was calibrated with a KG5 short pass-filtered silicon reference solar cell. The cells were then measured in both backward and forward direction with a constant scan rate of circa 0.6 V s⁻¹ (Keithley 2400 source measurement unit) while holding the temperature of the solar cell at 25 °C with a microcontroller-adjusted Peltier element. The MPP was tracked by using a perturb-and-observe method. HIF was calculated as fraction of PCE measured in backward and forward direction: HIF = PCE_BW/PCE_FW.

Time-resolved PL: For the time-resolved photoluminescence measurements, a self-made PL setup with a pulsed laser (532 nm, 5 kHz repetition rate, 800 ps pulse width, 0.5–2 nJ (respectively 100 nJ for the Supporting Information) pulse energy), an ACTON spectrometer, and a CCD camera (PIMAX512) with a gated mode was used, while keeping the samples in ambient atmosphere. For the evaluation of the PL measurements three main processes were assumed: (1) Recombination with trap states (Shockley–Read–Hall) on the timescale of microseconds, (2) band-to-band recombination on short time scales (<0.1 μs), and (3) quenching of PL by extraction of holes. A two-term exponential fit was used to determine a lifetime constant τ, which was dominated by the process (1) and a second time constant τ₂, here called mixed time constant τ, which was dominated by process (2) and (3): I = a₁ exp (−t/τ₁) + a₂ exp (−t/τ₂). It has to be mentioned that since process (2) and (3) cannot really be described throughout an exponential decay function, this is only a rough estimation for the short time scales.

X-ray Diffraction (XRD): The X-ray diffraction patterns were measured using a Bruker D2 PHASER (Cu K-α radiation).

XPS Mapping: Samples for XPS mapping experiments were prepared as described above and transferred to the ultrahigh vacuum chamber of the XPS system (Thermo Scientific ESCALAB 250Xi). XPS measurements were performed using a X86 monochromated Al K-α source (hv = 1486.6 eV) and a pass energy of 20 eV.

Surface Profiles and Thickness: The thickness of the perovskite films was measured using a Bruker Dektak XT profilometer.

Contact Angle, SFE, and SFT: For calculating the wetting envelope and therefore the polar and dispersive part of the SFE of the substrates, the contact angle of four solvents (deionized water, diiodo methane, dimethyl sulfoxide, and ethylene glycol) was measured using a sessile drop method (Krüss DSA 100 drop shape analyzer system). Droplets (~1–5 μl) were set onto the surface and then measured after a short settling time. The polar and dispersive SFE were calculated using Owens–Wendt–Rabel–Kaelble theory using a least absolute residual method. The dispersive and polar part of the SFT of the inks were calculated with the total SFT, measured with the pendant drop method, and the measured contact angles on a PTFE substrate.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

The financial support by the Federal Ministry for Research and Education (BMBF) through the project PRINTPERO (03SF0557A), the Bundesministerium für Wirtschaft und Technologie through the project CAPITANO (3EE10388), the Initiating and Networking Funding of the Helmholtz Association (HYIG of U.W.P. (VH-NG-1148)); Recruitment Initiative of B.S.R.; the Helmholtz Energy Materials Foundry (HEMF); PEROOSE (ZT-0024): the project HYPer as part of HeiKα research collaboration; and the Science and Technology of Nanostructures Research Program as well as the Karlsruhe School of Optics & Photonics (KSOP) is gratefully acknowledged. The authors would like to thank the members of the Perovskite Taskforce at LTI and the Printed Electronics research group at iL for fruitful discussions and support in their scientific work, especially Dr. Bahram Abdollahi and Pariya Nazari for their help in supporting measurements. In addition the authors would like to express their gratitude toward Dr. Guillaume Gomard for supporting measurements and Marius Jakoby for very helpful discussion on photoluminescence measurements.

Conflict of Interest

The authors declare no conflict of interest.

Author Contributions

H.E. and F.S. contributed equally to this work. The manuscript was written through contributions of all authors.

Keywords

high diffusion lengths, inkjet printing, large columnar crystal grains, perovskite solar cells

Received: September 27, 2019
Revised: November 19, 2019
Published online: December 19, 2019

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