Comparative Performance Study of Perovskite Solar Cell for Different Electron Transport Materials

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Abstract

In recent times, planar organo-metal halide perovskite solar cells (PSCs) achieved high power conversion efficiency (PCE > 22%). Mixed organic-inorganic halide perovskites, with excellent light harvesting properties, have evolved as a promising class of semiconductors for photovoltaics. In this work, compositional and electrical characterizations of materials used for different layers of PSC have been studied. One dimensional solar cell simulator wx-AMPS is used for numerical simulation of such devices and all simulations are done under AM1.5 illuminations and 300K temperature. Investigating the influences of thickness of electron transport material (ETM), hole transporting material (HTM) and absorber on the photovoltaic performance of PSCs, it is observed that, increase in thickness of perovskite (MAPbI3) results in increase in PCE of solar cells, whereas increase in thickness of ETM layer results in decrease in the efficiency of the devices. The ETM plays a vital role on the performance of PSC. In this paper, for the first time performances of PSC for three different ETMs (TiO2, ZnO or SnO2) are calculated and analyzed simultaneously with the simulator wx-AMPS. The photovoltaic performances have been explored and efficiencies of 27.6%, 27.5% and 28.0% are reported for perovskite solar cells with TiO2, ZnO and SnO2 as ETM respectively for a specific thickness. Finally, this simulation study concludes that ZnO and SnO2 may be effective alternatives of the commonly used material, TiO2 as they are economically more potential and give somewhat better photovoltaic performance.

Keywords: Perovskite Solar Cell, Electron Transport Material, wx-AMPS.

I. Introduction

During the last decade, a new kind of dye-sensitized solar cell (DSSC) has been developed, where a solid hole conducting material is used in place of the liquid electrolyte. Michael Gratzel et al. fabricated solid state DSSC’s, by depositing two different solutions of CH3NH3I and PbI2 that form a composite dye of perovskite and achieved 15.0% of efficiency.

Among existing photovoltaic technologies, perovskite solar cells have unfolded as a promising alternative because of its superior optoelectronic behavior, solution-processability and excellent photovoltaic performance. The perovskite acts as a good absorber for solar cells as it has fascinating electron and hole diffusion ability coupled with very wide absorption wavelength range and high photo-conversion efficiency. Bach et al. developed the PSC by replacing the liquid-electrolyte with solid-state hole-transporting material (HTM) 2,2’7,7’tetrakis (N,N-di-p-methoxyphenyl-amine) 9,9’ –Spirobifluorene (OMeTAD).

Fig. 1 shows the basic device structure of a liquid-state DSSC and solid-state perovskite solar cells (PSC). Typically, fluorine-doped tin oxide (FTO) serves as the conducting glass substrate for DSSC. The substrate is followed by a TiO2 film, a porous nano-crystalline semiconductor oxide. TiO2 is normally sensitized bi-pyridine ruthenium complexes (dye) which absorbs visible light. The de-oxidized dye is oxidized by a redox electrolyte layer such as iodide/tri-iodide couple. The cathode is coated with platinum to collect electrons and catalyze the redox couple regeneration reaction. Holes are transported to the external circuit as the electrolyte regenerates the oxidized dye and transports the electrons back to the dye molecules.

Traditionally, CH3NH3PbI3 (perovskite) has long been used as sensitizer particles in dye-sensitized solar cells (DSSC) with liquid-electrolyte. Due to solvent leakage and corrosion of liquid electrolyte, researchers have been concerned on replacing this liquid electrolyte with a solid hole transporter to create a fully solid-state electronic Dye-sensitized solar cells. Miyasaka et al introduced perovskite in photovoltaics in 2009 and observed efficiency 3.81% and 3.13% for MAPbI3 and MAPbBr3 based solar cells, respectively. To this date, efficiencies of DSSCs have reached to 22.1%. However, in terms of efficiencies, cadmium telluride solar cells (CdTe, 19.6%), copper indium gallium diselenide (CIGS, 21.7%), crystalline silicon (c-Si, 25%), and gallium arsenide (GaAs, 28.8%) still outperform the PSCs. A deeper understanding of the material properties of the individual layers is necessary to explore the possibilities of achieving higher efficiencies in PSCs.

Fig. 1. Perovskite solar cell (Solid State) to Dye-sensitized solar cell (Liquid-state).

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In 2013, Mingzhen Liu et al. prepared vapor-deposited FTO/compact TiO$_2$/CH$_3$NH$_3$PbI$_3$/spiro-OMeTAD/Ag cell and achieved 15.4% efficiency$^{13}$. In 2014, a planar heterojunction CH$_3$NH$_3$PbI$_3$ solar cell with TiO$_2$ (ETM) and spiro-OMeTAD (HTM) were simulated with AMPS-1D and obtained 20% of efficiency by Feng Liu et al.$^9$. Weijun Ke et al. used the process of thermal oxidation to grow a thin TiO$_2$ film on the perovskite and prepared a cell that gave substantially high efficiency of 15.07%$^{16}$. In 2015, Jing Wang et al.$^{17}$ reported perovskite solar cells with thin dense Mg-doped TiO$_2$ as hole-blocking layers (HBLs), outperforming solar cells with TiO$_2$:HBLs due to their higher open-circuit voltage, $V_{oc} = 1.08$ V and PCE of 12.28%. Yun-Jeong Kim et al.$^{18}$, applied CH$_3$NH$_3$PbI$_3$ (perovskite) to FTO/glass and Au/Ti/polyethylene terephthalate (PET) substrate at 100°C. Very recently, Ao Zhang et al. achieved 26.6% of efficiency from ZnO(ETM)/MAPbX$_3$(perovskite)/CuSCN (HTM) -based solar cell$^{19}$.

In this work, a comparative study has been done by replacing TiO$_2$ with ZnO and SnO$_2$ as ETM layer with perovskite and spiro-OMeTAD. One major challenge of the commercialization of perovskite technology is that its long-term stability under thermal and moisture stresses is somewhat lower than that of inorganic devices. This drawback can be minimized to some extent by increasing the solar cell power conversion efficiency and reducing the fabrication costs. In this paper, it has been shown that efficiency can be improved and cost can be reduced if TiO$_2$ (commonly used ETM with PSC) is replaced with ZnO or SnO$_2$ as these materials are very cheap and available in market.

II. Modeling of Perovskite Solar Cell

In this study a simple model for perovskite solar cell of p-i-n device structure is used. The cross-section of this model is shown in Fig.2. The back contact is made of gold. The spiro-OMeTAD layer acts as the HTM and is located on top of the back contact. A low p-type doped perovskite (CH$_3$NH$_3$PbI$_3$) absorber layer is sandwiched between the n-type material (TiO$_2$) and the p-type spiro-OMeTAD layer. Here, TiO$_2$ is used as ETM with Fluorine doped Tin Oxide (FTO) on top.

In solar cell, absorber is important to harvest the energy by controlling doping concentration and charge carrier mobility; surface properties of the ETM and HTM affect the charge carrier conduction between the front and back contacts$^2$. Efficient collection of the charge carriers depends on work function of the front contact material and rear metallization. And most importantly the choice of ETM plays a vital role on the design of high efficiency perovskite solar cell as the energy band alignment between absorber and ETM layer is a crucial factor for the efficiency improvement of PSCs$^2$. TiO$_2$ is a semiconducting material and is widely used as ETM or hole-blocking layer with perovskite cells due to its appropriate band gap, electron affinity, large surface area, abundance and process ability in different forms.

![Fig. 2](image-url)

**Fig. 2.** Cross-sectional schematic diagram of p-i-n perovskite solar cell solar cell.

Some other mesoporous conducting oxides, ZnO and SnO$_2$ may be good alternatives when used with MAPbI$_3$ absorber and spiro-OMeTAD, hole transporting material. Fig.3 shows two proposed models, where ETM layer, TiO$_2$ is replaced by ZnO and SnO$_2$. Photovoltaic performances for PSCs with these three ETMs are determined by varying the thicknesses of absorber layer, ETM layer and HTM layer. FTO is taken as the front contact material as it is widely used with perovskite cell and Au has been taken as the rare contact.

III. Numerical Simulation: wx AMPS

The optimized structure of a PSC has been numerically modeled and simulated by using wx-AMPS (Analysis of Microelectronic and Photonic Structures) solar cell simulation software which is an improved version of AMPS-1D software under AM 1.5 G solar irradiation and 300K temperature$^{20}$. At first, optimized metal as back contact for PSC is determined by varying work function. Then optimized cell is obtained by varying ETM, absorber and HTM layer thicknesses. Effects of the variation of different layer thicknesses on photovoltaic performances (open circuit voltage, short circuit current, fill factor and efficiency) are discussed. Three different PSCs have been modeled by replacing ETM layer. Finally comparative discussion on the performances of various models is given with optimized the cell structures.

wx-AMPS solves Poisson’s equation and the hole and electron continuity equations. It is one of the reliable computer aided program for modeling various single and multi-junction as well as Schottky barrier devices both in hetero- and homo-structures$^{21}$. The assumptions made on this model are: carrier populations obey Maxwell–Boltzmann statistics, the only recombination mechanism is the radiative band-to-band process (monomolecular),
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internal conversion efficiency makes it to 100% and photon recycling effect appears in the cell. The general form of Poisson’s equation is given by equation 1.

\[
\frac{d}{dx} \left( -\varepsilon(x) \frac{d\psi(x)}{dx} \right) = q[n(x) - n_i(x) + N_p^+(x) - N_n^-(x) + \frac{n_i(x) - n_0(x)}{p_i(x) - n_0(x)}] \tag{1}
\]

where, \(N_p^+, N_n^-\) are the concentrations of ionized donors and acceptors, \(n, p\) are the free electrons and holes concentrations, \(n_i, p_i\) are the concentrations of trapped electrons and holes, \(\varepsilon\) is the dielectric permittivity of semiconductor, \(\psi\) is the electrostatic potential, and \(q\) is the charge of the carrier.

Solving the continuity equations provides the transport characteristics of the device. The continuity equations in steady state conditions are:

\[
\frac{dJ_n}{dx} = \frac{1}{q} \left( R_n(x) - G(x) \right) \tag{2}
\]

\[
\frac{dJ_p}{dx} = \frac{1}{q} \left( G(x) - R_p(x) \right) \tag{3}
\]

where, \(R_n, R_p\) are electrons and holes recombination velocities for direct band-to-band and indirect transitions, \(J_n, J_p\) are electron and hole current density, and \(G\) is the optical generation rate which is expressed as a function of \(x\) is,

\[
G(x) = -\frac{d}{dx} \sum_l \Phi_{l}^{FOR}(\lambda_i) + \frac{d}{dx} \sum_l \Phi_{l}^{REV}(\lambda_i) \tag{4}
\]

where, \(\Phi_{l}^{FOR}\) is the photon flux of the incident light and \(\Phi_{l}^{REV}\) is the photon flux the light reflected from the back surface. Both these parameters are defined at a wavelength, \(\lambda\) of \(l\) at some point \(x\). \(\Phi_{l}^{FOR}\) and \(\Phi_{l}^{REV}\) depend on the absorption and reflectance coefficient of light in the forward and reverse direction, respectively.

Solving equations (1), (2), and (3) yields state variables \(\psi(x), n(x), p(x)\) as a function of depth, \(\psi(x), n(x), p(x)\), and \(\psi(x), n(x), p(x)\). These quantities fully define the system at every point \(x\) in the device. Analytical solutions are not possible in this case.

In wx-AMPS, numerical technique, Newton- Raphson method is used and following boundary conditions are applied on the given set of equations\(^3\):

\[
\begin{align*}
\psi(0) &= \psi_0 - V \\
\psi(L) &= 0 \\
J_p(0) &= -qS_{po}[n_0(0) - p(0)] \\
J_p(L) &= qS_{po}[n(L) - p_0(L)] \\
J_n(0) &= qS_{no}[n(0) - n_0(0)] \\
J_n(L) &= -qS_{no}[n_0(L) - n(L)]
\end{align*} \tag{5}
\]

Here, \(S_{po}\) and \(S_{no}\) represent the effective interface recombination speeds for holes and electrons respectively at \(x = 0\) and \(S_{pd}\) and \(S_{nd}\) are effective interface recombination speeds at \(x = L\).

Electrostatic potential and the quasi-Fermi level for holes and electrons as a function of depth, at every point in the solar cell, are calculated in the wx-AMPS from the solutions of the three coupled differential equations as described above. Once these parameters are determined, the results can be used to obtain carrier concentrations, electric fields and currents, and device parameters like the open-circuit voltage (Voc), short circuit current density (Jsc), fill-factor (FF), and the efficiency (\(\eta\))\(^2\).

**IV. Result Analysis**

**Variation of ETM Thickness**

Fig.4 shows variation of open circuit voltage (Voc), short circuit current density (Jsc) and photovoltaic conversion efficiency (PCE) with thickness of ETM layer for three different cells. For three models, 400 nm thicknesses are taken for both MAPbI\(_3\) (absorber) and spiro-OMeTAD (HTM) layer. The thickness of ETM layer is varied from 10 nm to 450 nm and gradual decrease of Voc, Jsc and PCE is observed. The overall performance of solar cell with SnO\(_2\) as ETM is higher than other two models for the entire thickness range. At lower thickness (10 nm to 180 nm), TiO\(_2\)-based model gives better performance than ZnO-based cell, but for thickness higher than 250 nm ZnO-base model shows better efficiency than TiO\(_2\). The observation showed that TiO\(_2\) is more responsive to sunlight than that of the other two electron collecting materials as it has higher absorption co-efficient and reflectance and less transmittance\(^2\). It can be easily perceived that the increase in thickness of ETM degrade the performance of the solar cells. This is mostly due to fractional absorption of incident light by ETMs and variation in series resistance of the device with increasing thickness of ETM layer. Although better photovoltaic performances are found at thickness less than 50 nm, considering the chemical and mechanical stability of the device, 90 nm thickness of ETM may give moderately good result for all the three models. Efficiencies of 27.6 %, 27.5 % and 28.02 % are found for TiO\(_2\)-, ZnO- and SnO\(_2\)-base cell respectively at 90 nm thickness of electron collecting material.
Variation of HTM Thickness

Fig. 5 depicts variation of open circuit voltage ($V_{oc}$), short circuit current density ($J_{sc}$) and photovoltaic conversion efficiency (PCE) with thickness of spiro-OMeTAD (HTM) layer for all three models. Results are obtained for 90 nm thickness of ETMs and 400 nm of MAPbI$_3$ (absorber).

Variation of Absorber Thickness
Beyond this thickness efficiency is increased very slowly in the rate of photoion of open circuit voltage ($V_{oc}$) with the amount of light reaching and absorption range of the absorber is an optimum thickness is chosen by balancing between its absorption length and diffusion lengths. Moderately higher efficiencies have been found at 400 nm thickness of absorber, 90 nm thickness of ETM and 400 nm of HTM for the three models. Efficiencies of 27.6%, 27.5% and 28.02% are found for TiO$_2$, ZnO- and SnO$_2$-base cell respectively.

Among the three models, SnO$_2$-based cell has provided the best performance for entire thickness range of ETM layer. SnO$_2$ is an attractive electron collecting material for perovskite solar cell with FTO contact to enhance the performance. Compared to TiO$_2$, SnO$_2$ has higher electron mobility and deeper conduction band. This could accelerate the transportation of charge carriers from perovskite to electron transport layer, and also reduce the carrier accumulation at the surface. On the other hand, although ZnO-based cell gives somewhat lower performance than TiO$_2$-base cell, considering its lower cost, availability and higher electron mobility, ZnO can also be a good alternative of TiO$_2$ for perovskite solar cell. Also appropriate synthesizing of absorber, best selection of contact materials, and increasing conductivity of ETMs uplift the performance as well as stability of device.

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