Hybrid electroluminescent devices composed of (In,Ga)N micro-LEDs and monolayers of transition metal dichalcogenides

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We demonstrate a novel electro-luminescence device in which GaN-based µ-LEDs are used to trigger the emission spectra of monolayers of transition metal dichalcogenides, which are deposited directly on µ-LED surface. A special µ-LED design enables the operation of our structures even in the limit of low temperatures. A device equipped with a selected WSe₂ monolayer flake is shown to act as a stand-alone, electrically-driven single-photon source.

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I. INTRODUCTION

Atomically thin semiconductors, such as monolayers of transition metal dichalcogenides (TMDs) and their homo- and hetero-structures, exhibit interesting optical properties. The list of prominent attributes of these two-dimensional (2D) structures includes: strong excitonic binding which favors room temperature capabilities and possibilities to manipulate the valley degree freedom and emergence of new class of quantum emitters. In parallel to the intense fundamental research, pertinent efforts are also undertaken to exploit the TMD semiconductors in designing and fabricating of new optoelectronic devices of possible practical use. The operation of such devices is desired to be electrically driven. Controlled doping of TMD materials, although essential for the operation of conventional electroluminescent p-n diodes, still remains a largely unresolved challenge.

Therefore, alternative methods are being developed, relying most often on tunneling mechanisms of carrier injection into the active 2D component of the device. In fact, the electrically driven emission from TMD semiconductors has been successfully demonstrated in a number of differently designed structures. However, the architecture of such structures implies rather complex manufacturing processes, and the performance of the reported devices remains not yet fully satisfying. This calls for further efforts to improve/optimize the already proposed device schemes and/or to search for optional solutions.

Taking advantage of the advanced technology of nitride we put here forward a new concept of a compact pseudo-electroluminescent device by integrating the TMD semiconductors into (In,Ga)N light-emitting diodes (LEDs). Our devices are composed of micro-(In,Ga)N-LEDs (µ-LEDs), on top of which we directly deposit the TMD layers: the electrically driven (In,Ga)N µ-LEDs serve as the excitation sources to generate the PL emission from TMDs. A special design of our (In,Ga)N diodes ensures their functionality at cryogenic temperatures, which are preferential conditions for generating light from TMD monolayers and an unavoidable requirement when tracing the quantum emission centres from these layers. Low-temperature operation of our electroluminescent devices comprised of different: MoS₂, MoSe₂, WSe₂ and WS₂ monolayers, is demonstrated. Importantly, the hybrid (In,Ga)N-µ-LED/WSe₂-monolayer device is also shown to function as an electrically driven source of single photons. While the fabrication of the present devices implies a laborious exfoliation manufacturing, we speculate that the industrial scalability of such devices could be approached in the future, taking into account the progress in the growth of TMDs by epitaxy methods and thus evoking a possibility to grow the TMDs directly on (In,Ga)N LEDs, particularly with molecular beam epitaxy technique.

The technology of nitride LEDs has undergone intense development in recent years. Relevant efforts have been made to respond to pertinent demands for low energy consumption, that is, to achieve high luminous efficacy of LEDs, but also to improve their brightness and contrast, their shock resistance, and degradation time. Although many nitride LEDs already find relevant applications in display technology, the new possibilities for nitride optoelectronics have been raised more recently by employing a scheme of a tunnel junction. Such an architecture of (In,Ga)N LEDs is also crucial for the present work. The buried TJ enables the nitride LEDs to operate at cryogenic temperatures. Equally important is that the use of the TJ allows us to replace the top p-type layer (of a conventional LED) with a highly conductive n-type layer. This enhances the current spreading and enables the n-type application of the side contacts to the device. Thus, TMD flakes can be directly deposited on the GaN surface, and the experiments can be carried out at low temperatures.
II. EXPERIMENTAL DETAILS

The epitaxial structure, schematically shown in Figure 1, was grown by plasma-assisted molecular beam epitaxy (PAMBE) on (0001) bulk GaN crystals. The PAMBE grown TJ (90 nm) and µ-LED structure that emits light at a wavelength of 450 nm was capped with the n-type GaN of 150 nm (Figure 1b). The emission area of µ-LEDs was defined by a photolithography mask and He⁺ implantation (100 keV, \( D = 1.5 \times 10^{14} \text{ cm}^{-2} \)) of the µ-LED structure. Ion energy and dose were adjusted to penetrate the region at the location of the TJ (Figure 1c). The vacancies produced during the ion implantation of the TJ deteriorate the p-type and n-type conductivity and significantly increase the resistance of the TJ in the areas outside the µ-LEDs. After the ion implantation, the regrowth of 200 nm of high conductive n-type GaN enhances the current spreading at the top of µ-LEDs and allows the creation of the side contacts to the device. After regrowth, the 150 \( \times \) 300 \( \mu \text{m} \) devices containing arrays of individual microdevices were separated by reactive-ion etching. Next, on the Ga-polar side of the samples, the TJ/Al/Ni/Au metal contact was deposited with a photolithographic mask, followed by lift-off. The same metal contacts were used for bottom contact on the nitrogen polar side.

The TMD monolayers have been placed on top of (In,Ga)N µ-LEDs (see Figure 1a) using deterministic transfer procedures. Monolayer flakes were mechanically exfoliated from bulk crystals with chemically pure back-grinding tape and transferred to the µ-LED surface using Gel-Pak DGL-X4 elastomeric films. A number of devices have been fabricated, comprising different, MoS₂, MoSe₂, WSe₂ and WS₂ monolayers.

The optical measurements were carried out in the two experimental configurations. Time-integrated studies were performed with the samples placed in the cold-finger cryostat. The emission of TMD monolayers has been measured in two different configurations. Primarily, the measurements have been carried out using the external laser excitation at 532 nm (2330 meV). Secondly, the luminescence was triggered from underneath with (In,Ga)N µ-LEDs.

The spectra composed of sharp emission lines characteristic of a WSe₂ monolayer have also been examined with photon correlation experiments. These experiments have been carried out in the Hanbury-Brown-Twiss configuration, under (In,Ga)N LED excitation, and required the device to be immersed in superfluid helium to minimize the undesirable effects of heating (when large currents are driven to trigger the (In,Ga)N LEDs).

III. RESULTS AND DISCUSSION

The optical microscopy images of four different, WSe₂, MoSe₂, WS₂ and MoS₂ monolayers deposited on top of individual (In,Ga)N µ-LEDs are illustrated in Figure 2a-d. The characteristic, low temperature (10 K) microluminescence spectra of these layers are shown, correspondingly, in Figure 2e-h. The spectra shown with blue traces were measured under excitation with the external laser (532 nm), whereas those shown in red have been activated by the underneath (In,Ga)N µ-LEDs. The photoluminescence signals measured under external excitation are typical (with respect to spectral shape and characteristic energies and intensities) of TMD monolayers, as often studied when they are directly deposited on conventional
Si/SiO₂ substrates. Characteristic spectra of our WSe₂, MoSe₂ and WS₂ monolayers are fairly well reproduced when they are excited with (In,Ga)N μ-LEDs. On the other hand, the apparent background signal (see dashed traces in Figure 2), which is due to an unfiltered low-energy tail of the (In,Ga)N emission, largely distorts the spectrum of the MoS₂ monolayer. This is, in fact, because of a rather weak luminescence efficiency of MoS₂ monolayers, as commonly reported when this particular TMD monolayer is not protected by hexagonal boron nitride.

Having established the overall functionality of our pseudo-electroluminescence devices, we now turn our attention to structures displaying sharp emission lines, possibly comprising the quantum emitter centers. Such centers preferentially appear in WSe₂ layers. Several μ-(In,Ga)N-LED/WSe₂-monolayer devices have been fabricated and tested. The irregular shape of WSe₂ flakes has been found to favour the observation of sharp emission lines. One may speculate this is due to the increased probability of the formation of the specific centers trapping the photoexcited carries at the edges of TMD flakes or induced by inhomogenous strain distribution.

The representative example of a device displaying sharp emission lines contains a narrow, elongated flake of a WSe₂ monolayer. The device has been mounted on a cold finger of the cryostat, cooled down to 10K for primarily tested under such conditions. After turning on the (In,Ga)N diodes, the emission spectra, in the spectral range 1590 – 1775 meV, were measured when scanning the surface of the WSe₂ flake and its surrounding, with a spatial resolution of 2 μm. A map of the intensity of the PL integrated over the spectral range 1625 – 1640 meV is shown in Figure 3a. This image illustrates well the shape of the flake, which extends over the active areas of two neighbouring (In,Ga)N μ-LEDs. The background
signal due to the emission tail of the (In,Ga)N diode is fairly weak - see the spectrum measured outside the flake area (Figure 3b). The spectra measured in the areas in the middle of the flake, away from its edges, exhibit a broad feature, centered at 1660 meV – a typical response of WSe2 ML with significant spatial inhomogeneity (see Figure 3b). However, near the edges and narrow parts of the flake, the luminescence spectra often display sharp lines; see Figure 3b.

Notably, the low activation current (below 1 nA µm⁻²) of our device permits the measurements of the emission spectra of the WSe₂ monolayer over a broad range of current densities driving the (In,Ga)N diodes and thus enables the straightforward tuning of the excitation power to match the saturation point of a particular emitter. To illustrate this effect, we present a typical PL intensity of the emitter as a function of the current density (Figure 3c). As the luminescence intensity of narrow lines in WSe₂ monolayers is known to drop down upon increase of temperature, we conclude that the observed effect is likely due to a low cooling efficiency of our device under the current arrangement (limited by the heat transfer effectiveness between our device and the mounting base in the cryostat). However, the heat generation process can be suppressed by reducing the size of the µ-LED and the precise positioning of the TMD flakes.

To demonstrate the operation of our device as a single-photon source, the device was immersed in a superfluid helium bath and cooled down to a temperature of 1.65 K. This resulted in an improved cooling efficiency of the device, enabling the application of higher diode currents, and at the same time gaining higher emission intensities of narrow emission lines. This allowed for the correlation measurements of the photon statistics emitted from the particular centre.

In Figure 3a, we present the spectrum of the narrow line emitter in the WSe₂ monolayer under the (In,Ga)N µ-LED excitation that was powered with 1.8 µA µm⁻² current. We note the particularly low ratio of the background µ-LED spectral tail to the narrow emission line intensity. Any uncorrelated photons detected in the Hanbury-Brown-Twiss experiment would proportionally reduce the amplitude of coincidence dip as well as the feasibility of our device as a single-photon source.

The acquired photon coincidence correlation function, $g_2(\tau)$, is presented in Figure 3b. The photon coincidence correlation function $g_2(\tau)$ denotes the relative probability density of an event with a time separation $\tau$. The decrease in the $g_2$ value around $\tau = 0$ evidences the single-photon component in the emission. The emission has predominantly single-photon characteristics when $g_2(0) < 0.5$.

The amplitude of the correlation function at zero delay is significantly low ($g_2(0) = 0.22$), which attests to the single-photon character of the emission. The photon coincidence correlation can also be used to determine the lifetime of the emitter state (Figure 3b). The lifetime is equal to the characteristic decay time of the antibunching feature in the $g_2(\tau)$ profile at the low photon rate limit. The lifetime obtained in our experiment ($\tau_{\text{decay}} = 4.8 \text{ ns}$) is in agreement with the results obtained in other studies.

We also used the photon coincidence correlation experiment to verify the temporal stability of the emitter. Fluctuations of either the emitter intensity or the intensity of its excitation source would result in the effect of photon bunching. We observe no decrease of $g_2$ at large $\tau$, confirming the good stability of the emission center and the µ-LED. We also observed good long-term stability of our device. During several days of measurements, we have observed no significant variation in the emission efficiency of the selected center.

![Figure 4](image_url)

**FIG. 4:** (a) Luminescence spectrum of the narrow emission line chosen for the single-photon correlation experiment. (b) Coincidence correlation counts as a function of the time interval between photons. The photoluminescence is excited with µ-LED. The broad background luminescence from the µ-LED was not subtracted in the post-processing.

IV. CONCLUSION

In conclusion, we have demonstrated that light emission from TMD monolayers can be efficiently generated when they are directly deposited on, and excited by the (In,Ga)N µ-LEDs. Our hybrid devices, 100 µm in size, can be seen as an optional route towards the effective, electrically driven generation of light from TMD semiconductors. The devices equipped with different MoS₂, WSe₂, WS₂ and MoSe₂ monolayers have been shown to be operational at low temperatures thanks to the unique design of our (In,Ga)N micro-LEDs. A proof of concept for a new, electrically driven source of single photons has been demonstrated with one of our (In,Ga)N µ-LEDs/WSe₂-monolayer devices.
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1 Y. You, X.-X. Zhang, T. C. Berkelbach, M. S. Hybertsen, D. R. Reichman, and T. F. Heinz, Nature Physics 11, 477 (2015).
2 K. He, N. Kumar, L. Zhao, Z. Wang, K. F. Mak, H. Zhao, and J. Shan, Physical review letters 113, 026803 (2014).
3 J. S. Ross, S. Wu, H. Yu, N. J. Ghimire, A. M. Jones, G. Alavazian, J. Yan, D. G. Mandrus, D. Xiao, W. Yao, et al., Nature communications 4, 1474 (2013).
4 E. Lorchart, S. Azzini, T. Chervy, T. Taniguchi, K. Watanabe, T. W. Ebbesen, C. Genet, and S. Berciaud, ACS Photonics 5, 5047 (2018).
5 K. F. Mak, C. Lee, J. Hone, J. Shan, and T. F. Heinz, Physical review letters 105, 136805 (2010).
6 A. Splendiani, L. Sun, Y. Zhang, T. Li, J. Kim, C.-Y. Chim, G. Galli, and F. Wang, Nano letters 10, 1271 (2010).
7 H. Zeng, J. Dai, W. Yao, D. Xiao, and X. Cui, Nature nanotechnology 7, 490 (2012).
8 K. F. Mak, K. He, J. Shan, and T. F. Heinz, Nature nanotechnology 7, 494 (2012).
9 M. Koperski, K. Nogajewski, A. Arora, V. Cherkez, P. Mallet, J.-Y. Veuvillen, J. Marcus, P. Kossacki, and M. Potemski, Nature nanotechnology 10, 503 (2015).
10 P. Tonndorf, R. Schmidt, R. Schneider, J. Kern, M. Buscema, G. A. Steele, A. Castellanos-Gomez, H. S. van der Zant, S. M. de Vasconcellos, and R. Bratschitsch, Optica 2, 347 (2015).
11 M. Brooks and G. Burkard, Physical Review B 97, 195454 (2018).
12 J. Kern, I. Niehues, P. Tonndorf, R. Schmidt, D. Wigger, R. Schneider, T. Stiehm, S. Michaelis de Vasconcellos, D. E. Reiter, T. Kuhn, et al., Advanced Materials 28, 7101 (2016).
13 S. Kumar, A. Kaczmarczyk, and B. D. Gerardot, Nano letters 15, 7567 (2015).
14 A. Branny, S. Kumar, R. Proux, and B. D. Gerardot, Nature communications 8, 15053 (2017).
15 Y. Luo, N. Liu, X. Li, J. C. Hone, and S. Strauf, 2D Materials 6, 035017 (2019).
16 M. Blauth, M. Jürgensen, G. Vest, O. Hartwig, M. Prechtl, J. Cerne, J. J. Finley, and M. Kaniber, Nano Letters 18, 6812 (2018), ISSN 1530-6984.
17 F. Peyskens, C. Chakraborty, M. Muneeb, D. Van Thourhout, and D. Englund, Nature Communications 10, 4435 (2019), ISSN 2041-1723.
18 J. Wang, I. Verzhbitskii, and G. Eda, Advanced Materials 30, 1802687 (2018), ISSN 1521-4095.
19 H. Qiu, T. Xu, Z. Wang, W. Ren, H. Nan, Z. Ni, Q. Chen, S. Yuan, F. Miao, F. Song, et al., Nature Communications 4, 2642 (2013), ISSN 2041-1723.
20 S. Wang, W. Zhao, F. Giustiniano, and G. Eda, Physical Chemistry Chemical Physics 18, 4304 (2016), ISSN 1463-9084.
21 C. Palacios-Berraquero, M. Barbone, D. M. Kara, X. Chen, I. Goykhman, D. Yoon, A. K. Ott, J. Beitner, K. Watanabe, T. Taniguchi, et al., Nature Communications 7, 12978 (2016), ISSN 2041-1723.
22 R. S. Sundaram, M. Engel, A. Lombardo, R. Krupke, A. C. Ferrari, P. Avouris, and M. Steiner, Nano Letters 13, 1416 (2013), ISSN 1530-6984.
23 Y. Ye, Z. Ye, M. Gharghi, H. Zhu, M. Zhao, Y. Wang, X. Yin, and X. Zhang, Applied Physics Letters 104, 193508 (2014), ISSN 0003-6951.
24 R. Cheng, D. Li, H. Zhou, C. Wang, A. Yin, S. Jiang, Y. Liu, Y. Chen, Y. Huang, and X. Duan, Nano Letters 14, 5590 (2014), ISSN 1530-6984.
25 F. Withers, O. Del Pozo-Zamudio, A. Mischenko, A. P. Rooney, A. Golinia, K. Watanabe, T. Taniguchi, S. J. Haigh, A. K. Geim, A. I. Tartakovskii, et al., Nature Materials 14, 301 (2015), ISSN 1476-4660.
26 S. Wang, J. Wang, W. Zhao, F. Giustiniano, L. Chu, I. Verzhbitskii, J. Zhou, Yong, and G. Eda, Nano Letters 17, 5156 (2017), ISSN 1530-6984.
27 D. Li, R. Cheng, H. Zhou, C. Wang, A. Yin, Y. Chen, N. O. Weiss, Y. Huang, and X. Duan, Nature Communications 6, 7509 (2015), ISSN 2041-1723.
28 T. Wu, C.-W. Sher, Y. Lin, C.-F. Lee, S. Liang, Y. Lu, S.-W. Huang, W. Guo, H.-C. Kuo, and Z. Chen, Applied Sciences 8, 1557 (2018).
29 W. P cuski, M. Grzeszczyk, K. Nogajewski, A. Bogucki, K. Oreszczuk, J. Kucharek, K. E. Polczyńska, B. Seredyński, A. Rodek, R. Bożek, et al., Nano Letters 20, 3058 (2020), ISSN 1530-6984.
30 J. Day, J. Li, D. Y. C. Lie, C. Bradford, J. Y. Lin, and H. X. Jiang, Applied Physics Letters 99, 031116 (2011), ISSN 0003-6951.
31 Z. Chen, S. Yan, and C. Danesh, Journal of Physics D: Applied Physics 54, 123001 (2021), ISSN 0022-3727.
32 C. Skierbiszewski, H. Turski, M. Zak, K. Nowakowski-Szkudlarek, G. Muziol, M. Siekacz, A. Feduniewicz-Zmuda, and M. Sawicka, in 2018 76th Device Research Conference (DRC) (2018), pp. 1–2.
33 S. J. Kowser, E. C. Young, B. P. Yoonke, C. D. Pynn, R. M. Farrell, J. S. Speck, S. P. DenBaars, and S. Nakamura, Optics Express 25, 3841 (2017), ISSN 1094-4087.
34 C. Skierbiszewski, G. Muziol, K. Nowakowski-Szkudlarek, H. Turski, M. Siekacz, A. Feduniewicz-Zmuda, A. Nowakowska-Szkudlarek, M. Sawicka, and P. Perlin, Applied Physics Express 11, 034103 (2018), ISSN 1882-0786.
35 M. Chlipala, H. Turski, M. Siekacz, K. Pieniak, K. Nowakowski-Szkudlarek, T. Suski, and C. Skierbiszewski, Optics Express 28, 30299 (2020), ISSN 1094-4087.
36 K. Oreszczuk, T. Kazimierczuk, T. Smoleński, K. No-
gajewski, M. Grzeszczyk, A. Łopion, M. Potemski, and P. Kossacki, Physical Review B 102, 245409 (2020).

37 A. Srivastava, M. Sidler, A. V. Allain, D. S. Lembke, A. Kis, and A. Imamoğlu, Nature nanotechnology 10, 491 (2015).

38 Y.-M. He, G. Clark, J. R. Schaibley, Y. He, M.-C. Chen, Y.-J. Wei, X. Ding, Q. Zhang, W. Yao, X. Xu, et al., Nature nanotechnology 10, 497 (2015).

39 C. Chakraborty, L. Kinnischtzke, K. M. Goodfellow, R. Beams, and A. N. Vamivakas, Nature nanotechnology 10, 507 (2015).