Surface decorated Ni sites for superior photocatalytic hydrogen production

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Abstract
Precise construction of isolated reactive centers on semiconductors with well-controlled configurations affords a great opportunity to investigate the reaction mechanisms in the photocatalytic process and realize the targeted conversion of solar energy to steer the charge kinetics for hydrogen evolution. In the current research, we decorated isolated Ni atoms on the surface of CdS nanowires for efficient photocatalytic hydrogen production. X-ray absorption fine structure investigations clearly demonstrate the atomical dispersion of Ni sites on the surface of CdS nanowires. Experimental investigations reveal that the isolated Ni atoms not only perform well as the real reactive centers but also greatly accelerate the electron transfer via direct Ni–S coordination. Theoretical simulation further documents that the hydrogen adsorption process has also been enhanced over the semi-coordinated Ni centers through electronic coupling at the atomic scale.

KEYWORDS
electron transfer, hydrogen evolution, molecular activation, photocatalysis, single atom catalyst
1 | INTRODUCTION

The highly increased energy demand of modern society has attracted enormous attention for developing renewable and clean energy sources, among which solar energy is considered as a promising candidate.1–7 Photo-driven water splitting has been regarded as a sustainable and cost-efficient strategy to generate clean hydrogen (H₂) fuel to realize a stable supply of renewable energy.3,8,9 Semiconductor materials have received great research attention due to their unique chemical, physical, and optoelectronics characteristics. Nevertheless, the catalytic performance over those semiconductive materials is relatively low even under the participation of sacrificial agents. Recombination of photoexcited electron–hole pairs before the redox reaction is considered as the main reason for the low efficiency.10–12 Cocatalysts are usually applied in photocatalytic systems to accelerate the separation of photogenerated electron/hole and in-fine improving the reaction kinetics. Numerous researches have confirmed that Pt affords the highest hydrogen evolution rates among various cocatalysts. However, the practical utilization of Pt-based cocatalysts is highly hindered by their high cost and scarcity. Therefore, developing cost-efficient and highly active alternatives to Pt is urgently expected.13,14 The development of precious metal-free cocatalysts that can be readily used for low-cost and pilot system design is thus significantly favored but still challenging to date.

Generally speaking, charge separation through the capture of electrons by a cocatalyst affords the most powerful strategy to enhance the hydrogen evolution rates in the reaction solution.15–18 Manipulating the size of the cocatalyst down to a single-atom state represents the highest utilization of metal species, which may also improve the photogenerated electron–hole separation efficiency.19–23 Furthermore, homogeneously dispersed isolated reactive centers offer us an ideal model to explore the precise structure–performance correlation at the atomic level. The electronic structure of isolated reactive centers as cocatalysts and their interaction with light absorbers are important factors in the comprehension of the photocatalytic mechanism.24–27 Single-atom catalysts, with nearly identical catalytically reactive centers, have emerged as excellent candidates for photocatalysts, affording us the great chance to detect their structural evolution during the catalytic process.5 Moreover, the decoration of isolated active centers over semiconductors may also induce the generation of discrete energy bands, which is very effective in trapping charge carriers and boosting the transfer of photogenerated electrons and holes.28–30 Considering these characteristics, it is reasonable to expect greatly boosted catalytic performance for photocatalytic water splitting over single atom-decorated semiconductor photocatalysts.

Herein, we successfully decorated isolated Ni atoms (the cocatalyst) onto the surface of CdS nanowires (denoted CdS@Ni) for efficient photocatalytic H₂ evolution. Synchrotron radiation-based X-ray absorption spectroscopy has been applied to determine the coordination configuration of the isolated reactive centers. Photocatalytic water splitting performance evaluation undoubtedly demonstrates that the decorated Ni atoms work as the active centers for proton recombination, steering the charge kinetics and improving the photocatalytic activities dramatically. Density functional theory (DFT) simulations further demonstrate that the energy band configuration of the CdS semiconductor has been obviously changed after Ni decoration, resulting in an obviously reduced work function and greatly reduced Gibbs free energy for the adsorption of H*.  

2 | SYNTHESIS AND STRUCTURE CHARACTERIZATION

A solvothermal approach is applied for constructing CdS nanowires.8 The successful synthesis of CdS with a hexagonal wurtzite crystal structure is verified by X-ray diffraction (XRD) measurements (Figure S1). Scanning transmission electron microscopy (STEM) investigations verify the one-dimensional nanowires in morphology with a uniform dispersion in diameter (Figure S2). Isolated Ni atoms are further loaded onto the surface of the CdS nanowire via a facile wet-chemical method, followed by heat treatment to enhance the interaction between the isolated Ni atoms and CdS supports via the newly generated Ni–S bonding (Figure 1A). It should be noted that the crystal structure of the hexagonal wurtzite crystal structure of CdS did not change after the decoration of Ni species, as evidenced by the well-maintained XRD pattern, no peak indexed to the crystallographic Ni is presented (Figure S3). As confirmed by transmission electron microscopy (TEM) observation, no noticeable morphology of CdS nanowires can be detected with the decoration of isolated Ni atoms over the surface (Figure 1B,C). High-resolution TEM images with various magnifications of the CdS@Ni sample clearly reveal the absence of Ni or NiO₅ clusters (Figure 1D–F). The interplanar distance is confirmed as 3.4 Å, which can be ascribed to the (002) plane of hexagonal CdS, which is also consistent with the (002) spacing distance evaluated from the selected area electron diffraction spots (Figure 1G).31,32 The unvaried lattice fringes of CdS further demonstrate that the original structure of CdS is not changed after Ni decoration.33,34 All these results also suggest that the distribution of isolated
Ni atoms is either grafted on the CdS surface or should follow the periodic arrangement of Cd atoms in CdS.\textsuperscript{35,36} Catalysts with various Ni loading contents have also been synthesized. Inductively coupled plasma–optical emission spectrometry (ICP–OES) investigations confirm that the loading contents of Ni range from 0.29\% to 1.68\%. The Raman spectra have also been investigated for CdS- and Ni-decorated CdS@Ni, where the first-order LO Raman peak and second LO phonon vibrational peak appear at 297.9 and 598.1 cm\(^{-1}\), respectively (Figure S4).\textsuperscript{37} The observed LO Raman peak positions agree very well with those reported for CdS in the hexagonal wurtzite phase. The nearly identical vibration mode of these two samples also suggests that the Ni decoration cast limited influence on the surface of CdS. The Brunauer–Emmett–Teller surface areas of CdS and CdS@Ni were confirmed as 19.58 and 26.75 m\(^2\) g\(^{-1}\), respectively, indicating the highly dispersed states of the nanowires (Figure S5).\textsuperscript{38}

Extended X-ray absorption fine structure (EXAFS) spectrometry at the Ni K-edge has been explored using both wavelet transform (WT) and Fourier transformation to shed light on the local coordination geometry and oxidation state of the decorated Ni atoms. In the Fourier-transformed (FT)-EXAFS curve of CdS@Ni, a major peak around 1.86 Å assigned to the Ni–S coordination is clearly observed, while no Ni–Ni connection around 2.19 Å or Ni–O bond at 1.59 Å is detected, validating that Ni atoms are isolated and stabilized by the S atoms from the CdS without the formation of oxide and Ni aggregation (Figure 2A). The WT simulation has been conducted to afford radial distance resolution in the K space. The WT intensity maximum at about 6.1 Å\(^{-1}\) arising from the Ni–S coordination
is well resolved at 1.0–3.0 Å for CdS@Ni, whereas an intensity maximum at about 6.64 Å−1 associated with the Ni–Ni coordination is not detected (Figure 2E). These results further demonstrate that the Ni species in CdS@Ni should be isolated without any aggregation. The least-squares EXAFS curve-fitting analysis clarifies the coordination sphere of the atomically dispersed centers. The best-fitted results clearly confirm that the Ni–S bond length of 2.26 Å, much shorter than the Ni–Ni bonding (2.48 Å) in Ni foil (Figures 2B,C and S6). Further analysis demonstrates that the coordination number for the Ni–S bonding is about 4.3 in the first coordination sphere (Table S1).39,40 As shown in Figure 2D, the X-ray absorption near-edge structure (XANES) curve of CdS@Ni at the Ni K-edge has displayed varied near-edge absorption energy from that of Ni foil, suggesting that the Ni species exist with positive oxidation states.41,42 The absorption edges of XANES curves have been fitted, and the average oxidation state of Ni species in CdS@Ni is confirmed as 0.89 (Figure S7). Overall, the XAFS investigation clearly verifies that Ni species are atomically dispersed over the surface of the CdS nanowire via Ni–S coordination.

The successful decoration of Ni species over the surface of the CdS nanowire has also been verified by X-ray photoelectron spectroscopy (XPS) analysis (Figure 3A). The high-resolution Ni 2p spectrum in CdS@Ni with peaks at 879.3, 872.2, 863.4, and 854.8 eV are assigned to Ni 2p3/2, Ni 2p1/2 and their corresponding satellite peaks, which may suggest the generation of direct Ni–S coordination and exclude Ni–Ni bonding (Figure 3B).43–45 The dominant S 2p peaks located at 161.3 and 162.5 eV are attributed to lattice sulfur in CdS, while the minor peaks at 160.1 and 161.2 eV correspond to Ni–S bonds (Figure 3C).46,47 The binding energies of Cd 3d5/2 and Cd 3d3/2 are observed at 404.1 and 411.6 eV, revealing a normal Cd2+ state, which is ascribed to the Cd–S bonding existing in Ni–CdS (Figure 3D).48 In addition, compared with those of bare CdS, obvious shifts to a lower binding energy have been observed in both the Cd 3d and S 2p spectra, suggesting that the atomically dispersed Ni causes the increased electron density of the CdS support.49

3 | PHOTOCATALYTIC HYDROGEN EVOLUTION

After clarifying the atomic-scale structure of CdS@Ni, we further shifted our attention to the role of isolated Ni atoms in photocatalytic hydrogen evolution. Visible light (λ ≥ 400 nm)-driven photocatalytic H2 evolution is evaluated while lactic acid is applied as the sacrificial electron donor. The catalytic performance is also influenced by the volume ratio of lactic acid in the reaction system. The optimized performance can be obtained for CdS@Ni
when the volume ratio of lactic acid is 20% (Figure S8). Compared with the pure CdS nanowires, the decoration of Ni species tremendously boosted the catalytic performance. With the increase in the loading contents of Ni species, the H\textsubscript{2} evolution rates also increased accordingly (Figure 4A). Interestingly, the sample with a Ni loading content of 1.35% reached its maximum hydrogen evolution rate of 2.26 mmol g\textsuperscript{-1} h\textsuperscript{-1}/0.452 mmol h\textsuperscript{-1}, which is about 25 times higher than that of the bare CdS nanowires (0.09 mmol g\textsuperscript{-1} h\textsuperscript{-1}). This value is quite impressive and even comparable to those of precious metal-based single-atom photocatalysts (Table S2).\textsuperscript{25,50} We may thus be reasonable to conclude that the isolated Ni site performs as the real reactive centers in the hydrogen evolution process.\textsuperscript{32} Further increasing the loading content of Ni species has little influence on its catalytic performance, suggesting that the latter Ni sites are not involved in the rate-determining step above a certain surface concentration on CdS. Based on these facts, the optimal CdS@Ni with a Ni loading of 1.35 wt% is chosen for further photocatalytic studies. Control experiments were also conducted to clarify the nature of the hydrogen evolution process under photoirradiation (Figure S9). No H\textsubscript{2} can be detected in the absence of a sacrificial agent in the reaction system, confirming that lactic acid is indispensable in the hydrogen evolution process. The whole reaction was also immediately terminated without light irradiation. These results reveal that hydrogen evolution is indeed driven by photoirradiation.

The stability of the catalyst was further investigated via recycle testing and structural characterization after catalysis (Figure 4B). No obvious degradation of the catalytic performance is observed after five runs. XRD characterization of the sample after catalysis was conducted, and no change was monitored after the reaction, suggesting the high stability of CdS@Ni (Figure S10). The EXAFS and XANES spectra for the sample CdS@Ni after catalysis were also explored, and no obvious change was observed, suggesting that there is no Ni aggregation in the catalytic process and that the coordination configuration of isolated Ni atoms is well maintained after the reaction (Figure S11). The H\textsubscript{2} generation rates of CdS@Ni at different light wavelengths were also explored, in which the catalytic performance qualitatively tracks the visible light absorption spectra (Figure 4C). This result further demonstrates that hydrogen evolution is driven by photoexcitation, while the isolated Ni sites should be the active center for the outstanding photocatalytic activities of CdS@Ni.
FIGURE 4 (A) H₂ evolution over photocatalysts with different loading contents. (B) Cycling test for H₂ evolution over CdS@Ni. (C) Wavelength-dependent H₂ evolution rates and light absorption spectrum of CdS@Ni. (D) Steady-state photoluminescence (PL) spectra. (E) Time-resolved PL spectra. (F) Transient photocurrent responses of CdS@Ni and CdS

4 | MECHANISM EXPLORATION

The photogenerated electron/hole separation and transfer dynamics have also been conducted to clarify the origin of the greatly boosted photocatalytic performance (Figure S12). As observed, the intrinsic carrier recombination is quite obvious, as evidenced by the strong photoluminescence (PL) intensity (Figure 4D). After Ni decoration, the emission intensity of CdS@Ni is significantly damped, suggesting that the isolated Ni sites can improve the charge carrier separation with impressive trapping ability.

To further verify this tentative conclusion, the fluorescence decay times for CdS@Ni and CdS were also explored as 1.158 and 8.175 ns, respectively (Table S3). The significantly quenched fluorescence intensity and the reduced exciton lifetime clearly suggest a significantly accelerated migration of photoexcited electrons to the catalytic active Ni centers of CdS@Ni (Figure 4E). The energy transfer (ET) processes have also been quantitatively described via the measurement of the ET efficiency (ΦET) based on the time-resolved fluorescence decay lifetimes (Equation S1). The ΦET of CdS@Ni (86.7%) was confirmed based on donor lifetimes in the presence and absence of acceptor molecules, indicating that highly efficient electron transfer was realized after the decoration of Ni.

Transient photocurrent responses have been investigated to account for the charge-transfer process (Figure 4F). As expected, the photocurrent density of CdS@Ni is much higher than that of CdS, indicating that the Ni decoration exhibits a better capability to capture electrons and suppress the recombination of photoexcited electron–hole. Therefore, the isolated Ni atom could accelerate the surface reaction kinetics in the catalytic process. This viewpoint is further demonstrated by photoelectrochemical impedance spectroscopy (Figure S13). Lower charge-transfer resistance for CdS@Ni has been confirmed by the evidence that a smaller arc is detected. Photo-irradiation can further decrease the arc diameter, suggesting that the isolated Ni atoms act as electron collectors to facilitate the separation of photogenerated electron–hole pairs on the CdS@Ni interface. Thus, an adequate supply of electrons is realized by isolated Ni atoms, optimizing the elementary process of photocatalytic hydrogen evolution.

DFT simulations are also performed to understand the mechanism of the greatly enhanced performance of CdS@Ni. The electron rearrangement between the Ni atoms and S atoms from the CdS supports has been observed, as evidenced by the significantly modified differential electron densities (Figures 5A and S14). Compared with pure CdS, the Ni decoration also induces electron transfer from Ni atoms to S atoms and the generation of a new band in the vicinity of the Fermi level. These occupied states of Ni are evolved into the electron trap state for CdS and lift the Fermi level to a position close to the
conduction band minimum of CdS, indicating that Ni decoration can act as a mediator to promote electron transfer from CdS to isolated reactive centers (Figure 5B,C).\textsuperscript{59–61} Strong electronic coupling between the support and isolated Ni species is thus observed, which also reveals that Ni decoration can boost electron transfer with a narrower bandgap.

The variation in the electrostatic potential of CdS has also been explored (Figure 5D). As documented, the difference between the vacuum level and the Fermi level should numerically equal the work function ($\Phi$). Thus, the work function of the pure CdS (111) facet is confirmed as 2.4 eV.\textsuperscript{62} The variation in the electrostatic potential of CdS@Ni has also been investigated, while its work function has been verified as 2.08 eV (Figure 5E). All these results undoubtedly reveal that the work function of CdS@Ni is slightly reduced after Ni decoration, resulting in highly promoted photogenerated electron transfer from CdS to isolated Ni sites.\textsuperscript{63} The Gibbs free energy of hydrogen adsorption ($\Delta G_{\mathrm{H}^\circ}$) has also been explored in current research (Figure 5F). The calculated $\Delta G_{\mathrm{H}^\circ}$ for pure CdS is confirmed as 2.34 eV, suggesting that hydrogen is not favorably bound to the pure CdS surface. The hydrogen adsorptions are more favorably supported after Ni decoration, as evidenced by the greatly reduced $\Delta G_{\mathrm{H}^\circ}$ (0.47 eV) over the Ni sites, which is essential for high catalytic performance.\textsuperscript{64,65}

## Conclusion

In summary, isolated Ni atoms are decorated onto the surface of CdS nanowires via the newly generated Ni–S bonding. Experimental investigation and theoretical studies confirm that strongly stabilized Ni single atoms trigger a continuous and reversible photocatalytic hydrogen evolution process. The photocatalytic hydrogen evolution activity of the optimal CdS@Ni is about 25 times higher than that of the bare CdS nanowires, affording direct evidence that the isolated Ni atoms perform as the real active sites for hydrogen evolution. DFT simulation demonstrates that new electronic states are generated after Ni decoration, which is beneficial for hydrogen adsorption in the hydrogen evolution process, resulting in significantly enhanced catalytic hydrogen evolution capability. Current research will be attractive for the development of earth-abundant reactive centers and inspire the exploration of other semiconductor-based catalysts for photocatalytic energy conversion.

## Characterization Procedure

XRD patterns were characterized by X-ray diffractometer with graphite monochromatized Cu–Kα radiation (Bruker, D2 Phaser). UV–vis absorption was characterized by
UV–vis spectroscopy (Shimadzu, UV-2450). PL spectra and fluorescence decay curves were measured using a fluorescent spectrophotometer (Edinburgh, FLS800). Fluorescence decay curves were measured using a fluorescent spectrophotometer (Edinburgh, FLS800). TEM of the samples was performed with a Titan Themis-Z microscope from Thermo Fisher Scientific by operating it at an accelerating voltage of 300 kV and with a beam current of 0.1–0.4 nA. Dark field imaging was performed by STEM coupled to a high-angle annular dark-field (HAADF) detector. The STEM–HAADF data were acquired with a convergence angle of 20.9 mrad and a HAADF inner angle of 49 mrad. Furthermore, an X-ray energy dispersive spectrometer (FEI SuperX, s=0.7 sR collection angle, 10 eV dispersion) was also utilized in conjunction with DF-STEM imaging to acquire STEM–EDS spectrum-imaging datasets (image size: 512 × 512 pixels, dwell time 6 μs). Surface chemical analysis was performed by XPS (ULVAC-PHI Inc., PHI Quantera SXM). The molecular weight of the cluster was identified by ESI-MS (Thermo Scientific, Exactive Plus). The Ni contents were detected by ICP–OES (Thermo IRIS Intrepid II XSP spectrometer) after dissolving the samples in a mixture of HCl and HNO3 (3/1 in volume ratio).

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CONFLICT OF INTEREST
The authors declare no conflicts of interest.

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**SUPPORTING INFORMATION**

Additional supporting information can be found online in the Supporting Information section at the end of this article.

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