Hydroelasticity effects on water-structure impacts

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Received: 7 October 2019 / Revised: 24 July 2020 / Accepted: 26 July 2020 / Published online: 9 August 2020 © The Author(s) 2020

Abstract

Local and global loadings, which may cause the local damage and/or global failure and collapse of offshore structures and ships, are experimentally investigated in this study. The research question is how the elasticity of the structural section affects loading during severe environmental conditions. Two different experiments were undertaken in this study to try to answer this question: (i) vertical slamming impacts of a square flat plate, which represents a plate section of the bottom or bow of a ship structure, onto water surface with zero degree deadrise angle; (ii) wave impacts on a truncated vertical wall in water, where the wall represents a plate section of a hull. The plate and wall are constructed such that they can be either rigid or elastic by virtue of a specially designed spring system. The experiments were carried out in the University of Plymouth’s COAST Laboratory. For the cases considered here, elasticity of the impact plate and/or wall has an effect on the slamming and wave impact loads. Here the slamming impact loads (both pressure and force) were considerably reduced for the elastic plate compared to the rigid one, though only at high impact velocities. The total impact force on the elastic wall was found to reduce for the high aeration, flip-through and slightly breaking wave impacts. However, the impact pressure decreased on the elastic wall only under flip-through wave impact. Due to the elasticity of the plates, the impulse of the first positive phase of pressure and force decreases significantly for the vertical slamming impact tests. This significant effect of hydroelasticity is also found for the total force impulse on the vertical wall under wave impacts.

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1 Introduction

Hydroelasticity in marine applications is discussed in the early works of Chuang (1970), Bishop and Price (1979), Faltinsen (1997, 2000) and Faltinsen et al. (2004). The research by Faltinsen (1997, 2000) demonstrates theoretical and experimental results that are significant for the design of offshore structures: that the maximum bending stress of a structural plate section is proportional to the drop velocity and is neither sensitive to where the waves hit the wet deck nor to the curvature of the wave crest in the impact region (Faltinsen 1997, 2000). More recently, Temarel (2008) investigated the effects of hydroelasticity for a variety of marine structures such as ships, offshore platforms, very large floating structures and also aqua-culture structures, and the results show that hydroelasticity is important when predicting transient forces, such as slamming. There is a significant body of research on high-speed craft related the localized hydroelasticity effects, for
example how the slamming loads and structural responses are affected by the elasticity of the hull/panel (Aarsnes 1994; Kvålsvold 1994; Battley et al. 2009; Stenius 2009; Stenius et al. 2011a,b; Stenius et al. 2013). With increasing ship size, craft speed, and severity of environmental loadings, the localized hydroelasticity effects have become more of an issue for achieving optimized structures. In general, hydroelasticity can be considered as a sub-category of fluid–structure interaction between flexible structure and liquid. A classic example being the hull-water impacts of high-speed craft which can produce large transient hydrodynamic impact loads on the hull/bottom structure (von Karman 1929). There is a critical influence of the impact velocity and the relative angle between the hull and water surface (the deadrise angle) on the impact loads: the impact load increases with increasing impact velocity and decreasing deadrise angle (Wagner 1932; Zhu 1995). However, with small deadrise angles the impact pressure will be decreased by an air-cushioning effect (Verhagen 1967; Lewison and Maclean 1968; Faltinsen 2014). A flexible structure will be deformed under hydrodynamic loading and this deformation of the structure will affect the local flow-field between the structure and water, and thereby the spatial and temporal pressure distributions on the structure. Kinematic and inertia effects have been identified as two types of hydroelasticity effects during an impact event (Stenius et al. 2013); kinematic effects are associated with the structure response, i.e. the structural deformation changes the geometry, velocity and acceleration conditions at the fluid–structure boundary; and on the other hand inertia effects are associated with the rise time of loading of the structure. Kinematic and inertia effects are fundamentally combined and interrelated, however a distinction between them has been discussed by Stenius (2009) and Stenius et al. (2013). Kimmoun et al. (2009) have investigated hydroelasticity experimentally by considering wave impacts on a flexible vertical wall. Their study investigated pressure distribution on the flexible wall, and deflection of the wall under various types of impact. Further, the kinematics of the fluid and fine details of the flow (air pocket) were investigated experimentally and theoretically. Their study showed fairly good agreement between these approaches.

In this present study, hydroelasticity effects on impact were experimentally investigated in a novel way by dropping a flat plate from various heights onto the water surface, with zero deadrise angle. A spring system was connected between the impact plate and the carriage to form an elastic structure, to understand the role of elasticity effect on slamming impacts on ship and offshore structures. In addition, to gain a better understanding of the physical processes involved in breaking wave impacts on a large ship or offshore structure (where elasticity of structure may have an important effect on wave impacts), various types of wave impact on a truncated vertical wall, considered to be part of a FPSO hull, have been experimentally investigated in this work. The truncated wall is an initially rigid wall that can be modified to an elastic wall by using a spring system on the rear side of the wall. Deflection of the springs, pressures and forces have been measured under the impact for both configurations.

2 Methodology

2.1 Drop test

The first series of experiments is designed to investigate hydroelasticity effects on slamming. It involves bodies slamming onto water, examples of which are the bottom and bow of a high-speed vessel, ditching of an aircraft, impact of bottom and bow of large ships like FPSO. It will model these interactions by performing a drop test with zero-degree deadrise angle.

The experimental work was carried out in the Ocean Basin in the University of Plymouth’s COAST Laboratory. The ocean basin is 35 m long by 15.5 m wide and has a movable floor that allows operation at different water depths up to 3 m. For these tests, the depth was set to 1 m. The falling block included a rigid impact plate connected to a carriage constructed from two driver plates with a total mass of 52 kg. A spring system was used to form elastic plates; elasticity of the plate could be changed using different spring stiffness (Mai et al. 2014; Mai 2017; Mai et al. 2019a). The falling block could be freely dropped in a 4 m high vertical guide frame fixed on the gantry crossing over the ocean basin. After each test, the falling block was lifted up to the tested drop height by a crane (see Fig. 1a, b). The impact plate was 0.25 m long, 0.25 m wide and 0.012 m thick. The impact velocity varied between 1 and 7 m/s by varying the drop height of the impact plate in the experiments to investigate the relationship between impact velocity, maximum pressure and force at impact. Pressures under the impact plate were measured by five miniature pressure transducers (FGP Sensors XPM10 having measurement range of up to 100 bar) installed at various locations on the lower face of the plate. The velocity of impact was estimated by integrating data recorded by an accelerometer (Model 4610 with range of up to 200 g or 500 g, in which g is the gravity acceleration) mounted on the top of the impact plate. A displacement sensor (Model LVDT-GCA500) was used to measure deflection of the springs under slamming impact. The geometry of the impact plate, the vertical distribution of mass and the configuration of the instrumentation on the impact plate are presented in Figs. 1c–e. 2. The sampling frequency was 50 kHz for pressure transducers, accelerometer and displacement sensors. This rate was sufficient to distinguish all the features of interest in the pressure and acceleration time histories. The velocity was integrated from the measured acceleration and the impact velocity is
the maximum velocity obtained at the time when the impact plate is about to hit the water surface. The impact velocity was varied from 1 to 7 m/s by changing the drop height.

### 2.2 Wave impact test

The second series of experiments was the offshore breaking wave impacts on a truncated vertical wall, representing a hull section of an FPSO. Various types of wave impact were generated and tested to identify the most violent impact type for ship and offshore structures. Furthermore, elasticity of the wall was tested to investigate how it affects impact loadings on the structure.

The experimental work was carried out in the sediment wave flume, also in the COAST Laboratory. The wave flume is 35 m long with a width of 0.6 m, a height of 1.2 m, and with a maximum still water depth of 0.8 m. A schematic of the physical model setup is given in Fig. 3. The truncated vertical wall (Plate 1) is an aluminium plate of 0.56 m width by 0.6 m height and 0.012 m thickness connected to a rigid part (Plates 2 and 3) by four springs. Plate 2 and 3 were mounted on a support frame via a low profile load cell and Plate 4. There were 0.02 m gaps on both sides of the tested model to ensure no friction between the model and the flume side walls, which would have affected the deformation of the springs and the load measurement. The spring system can incorporate springs of different stiffness and can also be locked to obtain a rigid wall impact model (Mai et al. 2015; Mai 2017; Mai et al. 2019a). Pressures under wave impact were measured by FGP XPM10 pressure sensors installed at 7 locations on the impact wall. A low profile load cell was used with an inline DC amplifier (Model 140) to measure total force on the wall. Accelerometers (Model 4610) were used to measure vibration of the structure under wave impact. The configuration of the instrumentation on the impact plate is presented in Fig. 3c. Figure 4 shows more details of the mass horizontal distribution of the tested walls. The pressure, force, deflection and acceleration data were sampled at 35 kHz frequency. This sampling frequency was smaller than that used in the drop test due to the larger number of measurement instruments and the RAM (Random-access memory) limitation of the computer used in this experiment.

Thirteen resistance-type wave gauges were used to measure water elevation along the wave flume, of which five wave gauges were used to measure wave profile in front of the
model during impact (Fig. 3b). The wave gauge data were sampled at 128 Hz frequency.

A focusing technique was applied to generate different types of wave impact by changing the focus location $X_f$ from the wave paddle (Kimmoun et al. 2010). Focus wave groups were generated using NewWave (Tromans et al. 1991) focusing with an underlying JONSWAP spectrum ($\gamma = 3.3$). Four different wave impact types were generated by changing the focus location. The distance between the front impact plate (Plate 1) and the wave paddle was 26.9 m. The target focus points were located at $X_f = 28.54$ m, 28.84 m, 29.04 m and 29.44 m from the wave paddle for the broken (Fig. 5a), high-aeration (Fig. 5b), flip-through (Fig. 5c) and slightly breaking (Fig. 5d) waves, respectively. The experimented wave characteristics had the significant wave height $H_s = 0.163$ m, peak wave period $T_p = 1.601$ s and the number of wave components $N = 116$. The wave conditions ($H_s$ and $T_p$) were scaled from prototype by a factor of 1:65 of the 100 year extreme significant wave height at the Cleeton platform in the Southern North Sea (Williams 2008). The tested water depth used in this study was 0.7 m. See Mai et al. (2015), Mai (2017) and Mai et al. (2019b) for more detailed information.

3 Results and discussions

3.1 Hydroelasticity effects on slamming

3.1.1 Typical time history of slamming

Typical motion time histories for the impact of the 52 kg rigid and elastic plates with about 4 m/s impact velocity, are shown in
Fig. 6a–c for the rigid plate, Fig. 6d–f for elastic plate 1 (using springs CXF51 × 54) and Fig. 6g–i for elastic plate 2 (using springs CXF51 × 102). The acceleration signals are clipped in Fig. 6a, d, g because the measured accelerations exceeded the 200 g and 500 g limits of the accelerometer in those tests. Figure 6b, e, h show that the velocity is zero at the start of the tests and then increases linearly during the free falling of the impact plates until a maximum velocity is obtained. At the instant of reaching its maximum velocity, the impact plates start to decelerate and this corresponds to the first contact with the water surface. The velocity of the impact plates fluctuates for about 30 ms, then decreases smoothly to zero and changes its direction due to the safety rope, which was connected to the carriage and used to stop the carriage from leaving the guide frame. The clipping of acceleration signals has an effect on the integrated velocity after the plate hit the water surface, but it does not affect the maximum velocity (the impact velocity). It is shown that the oscillations of acceleration and velocity increase with increasing elasticity of the plate (see small graphs in Fig. 6a, b, d, e, g, h. In Fig. 6c, f, i the integrated displacement of the impact plates during the test is shown.

The associated pressures and deflection of springs and the integrated force are presented in Fig. 7 for the rigid and elastic plates 1 and 2. Figure 7a, c, f present the pressures measured at locations P1 and P2 on the rigid plate, elastic plates 1 and 2, respectively (see Fig. 2 for the measured pressure locations). As was found for impact of the rigid plate (Fig. 7a), the impact pressures at locations near the edge (P2) are always attained earlier than the others at the centre of the impact plate (P1) and they are always much smaller than those at the plate centre under violent impacts (Fig. 7c, f). The impact pressure at the centre (P1) of the rigid plate is 14.75 bar (Fig. 7a) and this is higher than those on the elastic plate 1 (13.1 bar in Fig. 7c) and elastic plate 2 (12.73 bar in Fig. 7f). Similarly, the impact pressure near the edge (P2) of the elastic plate 1 (5.11 bar) and elastic plate 2 (5.05 bar) are slightly smaller than that of the rigid plate (5.38 bar).

Fig. 7 presents the integrated force and the measured deflection of springs. The maximum force of 37.82 kN on the elastic plate 1 due to impact with \( v = 4.11 \) m/s is very slightly higher than that on the elastic plate 2. The force on the elastic plate 2 is 37.81 kN due to impact with \( v = 4.11 \) m/s.
plate 2 due to impact with \( v = 3.90 \text{ m/s} \), where the maximum force was 37.00 kN. In contrast, the maximum deflection of the springs of the elastic plate 1 (\( D_{\text{max}} = 1.47 \text{ mm} \)) is proportionally much smaller than that for the elastic plate 2 (\( D_{\text{max}} = 1.87 \text{ mm} \)). The maximum deflection was attained at about 2 – 3 ms after the impact (\( t = 0 \)). This delay was also found in the drop test on regular waves presented by Faltinsen (1997, 2000). The natural period of a spring can be estimated from \( T = 2\pi (m/k)^{0.5} \), where \( m \) is the mass of spring in kg and \( k \) is the spring rate in N/m. Therefore, the natural periods of the springs CXF51 × 54 (\( m_1 = 0.41 \text{ kg} \) and \( k_1 = 765,538 \text{ N/m for a single spring} \)) and CXF51 × 102 (\( m_2 = 0.66 \text{ kg} \) and \( k_2 = 397,400 \text{ N/m for a single spring} \)) are calculated to be \( T_1 = 4.6 \text{ ms} \) and \( T_2 = 8.1 \text{ ms} \), respectively. The rise time of the maximum deflections is about 2.2 ms for both spring types Fig. 7e, h. This rise time is about a half of the natural period of the spring CXF51 × 54 (\( T_1 = 4.6 \text{ ms} \)) and about one quarter of the natural period of the spring CXF51 × 54 (\( T_2 = 8.1 \text{ ms} \)). The natural frequencies of the four springs and the impact plate are about 99.5 Hz and 67.5 Hz for the elastic plate 1 (using springs CXF51 × 54) and the elastic plate 2 (using springs CXF51 × 102), respectively. See Table 1 in Appendix A for the detailed estimation of the dry natural frequency of the four springs and the impact plate.

Distinct post-impact pressure oscillations under impact of the elastic plates onto the water surface were observed in this experiment. Pressure P1, P2 and P6 are presented and compared with the impact of the rigid plate in water in Fig. 8, 9, 10, in which the impact velocities were 1.28 ± 0.07 m/s, 4.08 ± 0.17 m/s and 6.96 ± 0.18 m/s, respectively. In Fig. 8, the first pressure peak decreases with increasing elasticity of the plate. The oscillation after the impact is believed to be caused by the oscillation of the trapped air between the impact plate and water surface, which has been discussed in previous studies (Ma et al. 2016; Mai 2017; Mai et al. 2019a). The FFT spectra of the measured pressures are given in the right hand plots in the figures. Under higher impact velocities (\( v = 4.08 \pm 0.17 \text{ m/s} \) and \( 6.96 \pm 0.18 \text{ m/s} \)), there are significant differences in the post-impact oscillations between the rigid plate and elastic plates with higher frequency peaks evidence for the elastic plates (Figs. 9, 10). Significant differences between the pressure traces for rigid and flexible panels were also found from the controlled water slam test of composite hull panels presented by Batley et al. (2009), Stenius (2009) and Stenius et al. (2011a). Higher oscillation frequency of the pressure under impact of an elastic plate was also observed by Tenzer et al. (2015). In the present study, it is clearly seen for the elastic plates that there are higher order frequencies in the pressure oscillations under high impact velocities and most of these oscillation frequencies are higher than the natural frequency (~ 500 Hz) of the falling block, carriage and support frame set-up for the rigid plate test (Figs. 9, 10). As it was found in the impact of the rigid plate in water (Ma et al. 2016; Mai 2017; Mai et al. 2019a), the high peak frequencies (139.9 Hz and 319.7 Hz) of the post-impact pressures (Figs. 8b, d, f, 9b, d, f, 10b, d, f) may due to the repeated compression and expansion of the different diameter trapped air bubbles when the impact plate is about to hit the water surface. The trapped air in flat plate impacts has been visually found in previous studies (Mai et al. 2014; Mayer and Krechetnikov 2018). The 139.9 Hz and 319.7 Hz are far from the 500 Hz natural frequency of the falling block, carriage and support frame set-up for the rigid plate test. These frequency (139.9 Hz and 319.7 Hz) are also higher than the natural frequencies of the four springs.
and the impact plate which are about 99.5 Hz and 67.5 Hz for the elastic plates 1 and 2, respectively. As a result of the shock wave traveling through the air gap along the width of the rigid plate (Mai 2017; Mai et al. 2019a), a second pressure shock was observed at an early stage ($t = 0.6$ ms) under impact velocity of $6.96 \pm 0.18$ m/s (the solid lines in Fig. 10a), but this second pressure shock appears at later stages, at about $t = 3.2$ ms and 3.6 ms for the elastic plates 1 and 2 (the dashed and dotted lines in Fig. 10a) and these later appearances of the second pressure shocks seem to be caused by a combined effect of the plate elasticity and the trapped air under the impact. There is no evidence of this second pressure peak from the previous experimental drop test studies in which plates with large deadrise angles ($5^\circ$ and $10^\circ$) were used and so no air bubble could be trapped under the plate. Similar to the impact of the rigid plate (Ma et al. 2016; Mai 2017; Mai et al. 2019a), the impact of elastic plates have shown many high frequency peaks of the corresponding discrete impact pressure events and these high frequencies are slightly different between the rigid and elastic plates (Figs. 8b, d, f, 9b, d, f, 10b, d, f).

### 3.1.2 Impact pressure and force on rigid and elastic plates

Figure 11 presents the impact pressures and forces under the impacts of the rigid and elastic plates. The impact pressures are presented in Fig. 11a for the measured location at the centre of the plate (P1). The impact forces are shown

![Fig. 8](image1)  
**Fig. 8** Oscillations of pressures (left) and their spectra (right) due to impact of rigid and elastic plates in water ($v = 1.28 \pm 0.07$ m/s, $m = 52$ kg). Note the different vertical scales

![Fig. 9](image2)  
**Fig. 9** Oscillations of pressures (left) and their spectra (right) due to impact of rigid and elastic plates in water ($v = 4.08 \pm 0.15$ m/s, $m = 52$ kg). Note the different vertical scales

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in Fig. 11b for 3–10 repeats. In general, there is a slight reduction in impact pressure and force only at the higher velocities, beyond 5 m/s. This limited reduction may due to the stiffness of springs used to form the elastic plates 1 and 2. The results presented by Tenzer et al. (2015) also failed to show a clear trend of hydroelasticity effects on pressure peaks, with the body they tested also being relatively stiff. Stenius et al. (2011a, b) showed that measured pressure magnitudes increase at the centre of the panel width of a wedge with increasing flexibility of the tested body with this trend being similar to their numerical simulations. They found a change of pressure magnitudes as a result of the local change in deadrise angle and impact velocity for a flexible body.

3.1.3 Impulses (pressure and force) on rigid and elastic plates

Impulse values can be obtained by integrating force or pressure over the duration of the impact. In the present study, the first impulse ($I_{\text{first}}$) of the impact is defined as the impulse of the first positive phase of the impact, i.e. the area A1 in Fig. 12 and the total impulse ($I_{\text{total}}$) is integrated from the start of the impact until the signal falls back to the noise level, i.e. the sum of the areas A1 to A8 in Fig. 12. It is noted that, the “total impulse” for the drop tests is calculated for a duration from the start time of an impact until the time when the impact plate was restrained by the rope used to stop the plate falling onto the basin floor. Therefore, the “total impulse” used in this analysis does not consider the entire motion of the object brought to rest and therefore there

Fig. 10 Oscillations of pressures (left) and their spectra (right) due to impact of rigid and elastic plates in water ($v = 6.96 \pm 0.18$ m/s, $m = 52$ kg). Note the different vertical scales

Fig. 11 Impact pressure P1 and force of rigid and elastic plates ($m = 52$ kg)
may be a difference between the total impulses for various impacts (Mai 2017; Mai et al. 2019a).

$I_{\text{p,first}}$ and $I_{\text{f,first}}$ of the rigid and elastic plates impact are presented in Fig. 13. The masses of the plates are all of 52 kg. It is shown that $I_{\text{p,first}}$ and $I_{\text{f,first}}$ of the rigid plate are generally higher than those for the elastic plates. This can be seen most clearly at the high impact velocities beyond 4 m/s. However, the total impulses of pressure and force ($I_{\text{p,total}}$ and $I_{\text{f,total}}$) presented in Fig. 14 show there is no clear difference between the total impulses of the rigid and elastic plates.

3.2 Hydroelasticity effects on wave impact on a vertical hull

3.2.1 Comparison of acceleration, force, deflection and pressures on the rigid and elastic walls

Typical time histories of acceleration, force, pressure, and spring deflection on the rigid and elastic walls are presented in Figs. 15, 16 and in Appendix B. The corresponding FFT spectra are also presented in those figures. It can be clearly seen that there is a low oscillation frequency (~ 37 Hz) after impact in the time histories and FFT spectra of acceleration, force and spring deflection of the tested walls (Figs. 15, 16 and in Appendix B). Evidence of this low frequency can also be seen in the pressure signals under high aeration and flip-through wave.
impacts on the rigid and elastic walls (Figs. 16 and 24). Using the theoretical natural frequency of air bubbles in water as derived by Minnaert (1933) and Hattori et al. (1994), the 37 Hz frequency of oscillation observed here corresponds to air-pocket and bubble radius of 88.2 mm. As visualised observation of the impacts by high speed camera (Mai et al. 2019b), there are no large air-pockets or bubbles with radii of about 88.2 mm. Therefore, it cannot be said that the 37 Hz frequency is due to air-pocket and bubbles under the impacts. This relatively low frequency is far from the dry natural frequencies of 25.3 Hz and 14.4 Hz for the elastic walls 1 and 2, respectively, but this

Fig. 15 Time histories (left) and their spectra (right) of acceleration, force and deflection due to flip-through impact

Fig. 16 Pressure time histories (left) and their spectra (right) due to flip-through wave impact on the walls. Note the different vertical scales
37 Hz frequency is identical to the 40 Hz natural frequency of the whole structure including the support frame, impact wall, plates 2–4 and four springs. The detailed estimation of dry natural frequency of the four springs and the impact wall is presented in Table 2 in Appendix A. Acceleration of the impact walls increases with increasing flexibility of the wall and there are higher oscillation frequencies for the more flexible walls (Figs. 15, 22, 24). However, for the slightly breaking walls by comparing single test cases between the tested assess how hydroelasticity affects the pressure on the very sensitive under wave impacts and it is difficult to 23e, 25e). Pressures on the tested walls are found to be of spring under the tested wave impacts (Figs. 15e, 21e, 25a, b). Maximum force of the wall also decreases with increasing flexibility of the wall, except for the slightly breaking wave impact whereas the maximum force on the elastic wall 1 is higher than that on the rigid wall (Fig. 25c). This reduction of the maximum force can be clearly seen for the high aeration and flip-through wave impacts (see Figs. 15c, 23c). As expected, deflection of springs increases significantly with decreasing stiffness of spring under the tested wave impacts (Figs. 15e, 21e, 23e, 25e). Pressures on the tested walls are found to be very sensitive under wave impacts and it is difficult to assess how hydroelasticity affects the pressure on the walls by comparing single test cases between the tested walls (Figs. 16, 22, 24). However, for the slightly breaking wave, which is known to be more quasi-hydrostatic than the other wave impact types, as also found for the force on the wall, the pressures on the elastic wall 1 are higher than those on the rigid wall, although they are seen to decrease with the more flexible wall (Elastic wall 2), see Fig. 26. The respective FFT spectra are presented in the right hand plots in Figs. 16, 22, 24, 26 for the time histories of pressures presented in the left hand plots. The spectra are plotted on a logarithmic scale for the y-axis to show more detail of the spectral tail form.

3.2.2 Impact pressure and impulse of wave impact

The maximum pressures at various levels (z/d = 0.071 to 0.39, in which d is the water depth and z is the positive upward level of pressure sensor referring to SWL) are presented in Fig. 17 for different types of wave impact on the rigid and elastic walls. For the tested elastic wall 2, pressures were measured at four levels of z/d = 0.071, 0.32, 0.36 and 0.39. The vertical axis is the dimensionless level z/d of the measured points on the wall and SWL is represented by z/d = 0. The horizontal axis is the logarithm of dimensionless impact pressure, P_max/(ρgd). In each plot of Fig. 17, the black circle represents the impact pressures on the rigid wall, while the red cross and blue square are respectively for the elastic walls 1 and 2. The solid, dashed and dotted lines connect the mean values of the measured maximum pressure at each level on the rigid wall, elastic walls 1 and 2, respectively. It can be seen from Fig. 17 that the distributions of the impact pressures on the tested walls are different under various wave impact types (broken, high aeration, flip-through and slightly breaking) and the highest impact pressures were attained at the levels above SWL which is similar to the finding by Hofland et al. (2011), while most previous works found the highest impact pressures occurred at and around SWL (Hattori et al. 1994; Oumeraci et al. 1993; Hull and Müller 2002; Bullock et al. 2007). This difference may be due to the use of focused wave groups here and by Hofland et al. (2011). At levels of z/d = 0.25 to 0.32, high impact pressures on the rigid and elastic walls were attained under broken wave impact (Fig. 17a). Under high aeration wave and flip-through impacts, the impact pressures are significantly higher further up the wall (greater z/d), except for the high aeration wave impact in which the medium value of the impact pressures at level z/d = 0.39 is slightly smaller than that at level z/d = 0.36 on both the rigid wall and the elastic wall 1 (Fig. 17b, c). For the high aeration case, the impact pressure appears to peak at level z/d = 0.36, while for the flip-through the impact pressure continues to increase with distance up the wall. The highest impact pressure on the rigid wall and the elastic wall 1 were 1.15 bar (P_max/(ρgd = 16.8) and 0.44 bar (P_max/(ρgd = 6.4) which were both recorded at level z/d = 0.39 under the flip-through impacts (Fig. 17c). On the other hand, the highest impact pressure of 0.56 bar (P_max/(ρgd = 8.2) was measured at z/d = 0.36 on the elastic wall 2 under high aeration wave impact (Fig. 17b). Figure 17d illustrates the distribution of the impact pressures on the rigid and elastic walls under

Fig. 17 Impact pressures (P_max/(ρgd)) on the walls
the slightly breaking wave impacts. The impact pressures at the high levels \((z/d = 0.25\) to 0.39\) are slightly larger than those at lower levels for all the tested walls. The impact pressures due to the slightly breaking wave impacts are much lower than those due to other wave impact types (broken, high aeration and flip-through). Considering single impact events, the highest impact pressures on the rigid wall \((P_{\text{max}} = 1.15 \text{ bar})\) are much higher than those on the elastic walls \((P_{\text{max}} = 0.44 \text{ bar for the elastic wall 1; } P_{\text{max}} = 0.56 \text{ bar for the elastic wall 2})\). However, the medium values of all recorded impact pressures are more or less the same for the rigid and elastic walls at all levels, except the impact pressures on the elastic wall 1 are much smaller than those on the rigid wall and the elastic wall 2 under the flip-through impact (Fig. 17c). For the broken, high aeration and flip-through wave impacts, large scatter of the impact pressures was found at high levels \((z/d = 0.25\) to 0.39\) on the walls (Fig. 17a, b, c). The slightly breaking wave impact caused the smallest scatter of impact pressures at all levels and the spatial variation of the maximum pressures on the wall is simply related to the local hydrostatic pressures on the walls (Fig. 17d). The large scatter of the impact pressures on a vertical wall under different wave impacts were also observed in previous studies (Hattori et al. 1994; Oumeraci et al. 1993; Hull and Müller 2002; Bullock et al. 2007).

The total pressure impulses \((I_P^{\text{total}})\) were obtained by integrating pressure values over the impact duration. Figure 18 presents the total pressure impulses at different levels on the rigid and elastic walls under the tested wave impact types. The black circle, red cross and blue square markers represent the total pressure impulses on the rigid wall, the elastic walls 1 and 2, respectively. The line joining mean values of the total pressure impulses is also presented in each plot of Fig. 18. It can be seen that the total pressure impulses on the rigid wall (the black solid lines) are approximately the same on the elastic wall 1 (the red dashed lines) at all levels due to the tested wave impact types. On the elastic wall 2, pressure impulses at levels from \(z/d = 0.071\) to 0.32 are much higher than those on the rigid wall and the elastic wall 1 due to all wave impact types (Fig. 18a–d). In general, the total pressure impulse decreases with increasing level on the walls and this is because of the submerged duration at lower levels being greater than those at higher levels even though the maximum pressure at the lower level is much smaller. This can be seen from the pressure time histories at different levels on the wall under wave impacts presented in the previous section (Fig. 16) and Appendix B (Figs. 22, 24, 26).

### 3.2.3 Impact force and impulse of wave impact

The mean impact force \((F_{\text{max}})\) and force impulse \((I_F)\) on the rigid and elastic walls due to different wave impacts are presented in Fig. 19. In each plot, the horizontal axis represents the impact force \((F_{\text{max}})\) on the elastic wall, while the vertical axes represent those on the rigid wall. The diagonal solid line is the 1:1 line.
impact, the impact force on the elastic wall 2 is higher than that on the rigid wall.

Figure 19b shows that as the elasticity of the wall increases, the mean total force impulse on the wall is larger for all tested impact types (the broken, high aeration, flip-through and slightly breaking impacts). The scatter in the total force impulse (Fig. 19b) is much larger than that of the impact force (Fig. 19a). It is found that the difference in total force impulse is about 10% of the lowest total force impulse for all investigated wave impact types.

3.2.4 Maximum deflection and impulse of springs

The mean of the maximum deflections of the springs used for the elastic walls is presented in Fig. 20a from various repeats presented in previous sections. Due to the high aeration and flip-through impacts, the average deflections of the springs of the elastic wall 1 are approximately the same as each other (~0.09 mm) and they are nearly double the average deflections due to the broken and slightly breaking wave impacts (~0.05 mm). The highest average deflection of springs of the elastic wall 2 is found for the high aeration wave impact (~0.3 mm). The average deflection due to the broken wave impact (~0.24 mm) is higher than that of the flip-through impact (~0.2 mm). Overall, the average deflections of the springs of the elastic wall 2 are much higher than those for the elastic wall 1. These are simply due to the spring rate of the elastic wall 2 (CL51×254: $k = 4\times37.7$ N/mm) is nearly one-third of the spring rate of the elastic wall 1 (CL51×102: $k = 4\times98.5$ N/mm). Because force is a function of deflection ($D$) and structure stiffness ($k$), then the deflection impulse can be another way to present the change in momentum under the impact. The total deflection impulses are presented in Fig. 20b, where we can see that the total deflection impulses of the elastic wall 2 are about ten times those of the elastic wall 1. Similar to the force impulse, it is interesting to find that there is very little difference between the deflection impulses (about less than 8% of the lowest deflection impulse) found in all investigated wave impact types.

4 Conclusions

This experimental study investigates two different types of hydrodynamic interaction to shed light on the role of elasticity in slamming and wave impacts. These can have local and/or global effects on offshore structures and ships, and may cause structural vibrations and local damage. Two different experiments were conducted: (i) free-falling of rigid and elastic flat plates onto water surface, which is designed to represent severe slamming impacts with zero degree deadrise angle; (ii) different types of wave impacts on truncated vertical rigid and elastic walls. The elasticity of the tested plates and/or walls has an effect on the hydrodynamic impact loadings for both slamming and wave impacts, though for the slamming impacts this was for high impact velocities only. The total impact force has been observed to decrease on the elastic walls under the high aeration, flip-through and slightly breaking wave impacts. However, the impact pressure only decreased under the flip-through wave impact. Furthermore, the elasticity also has an effect on the post-impact oscillations of the pressures for the slamming impacts, where the impulsive loadings of the first positive phase decreases with the elastic plates. Similarly, the total force impulse decreases with increasing elasticity of the wall for wave impacts, but the total pressure impulse on the wall increases for the more flexible wall. The effects of elasticity on hydrodynamic impact loadings need to be considered to assess the local/global loads which may cause local damage of a hull section or global structural response.

Acknowledgements This study is a part of the FROTH (Fundamentals and Reliability of Offshore Structures Hydrodynamics) project supported by the Engineering and Physical Science Research Council (EPSRC Grant EP/J012866/1). The FROTH project was led by the University of Plymouth and the collaborative partners include Oxford University, University of Bath, City University London and the Manchester Metropolitan University. The authors gratefully acknowledge the financial support provided by EPSRC and useful discussions with the project partners.

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Appendix A: estimation of natural frequency

See appendix Tables 1, 2.
Table 1  Estimation of dry natural frequency of the elastic plates 1 and 2 (drop test)

| Elastic plate | Spring type | Stiffness (single spring) (N/m) | Stiffness (four springs) (N/m) | Mass (single spring) (kg) | Mass (four springs) (kg) | T (single spring) (ms) | T (four springs) (ms) | Freq. of single spring (Hz) | Mass of impact plate (kg) | Mass of four springs + impact plate (kg) | Freq. of four springs + impact plate (Hz) |
|---------------|-------------|--------------------------------|--------------------------------|--------------------------|-------------------------|------------------------|------------------------|---------------------------|------------------------------|------------------------------------------|-------------------------------|
| Elastic plate 1 | CXF51×64    | 765,538                        | 3,062,150 0.41                | 1.64                     | 4.60                    | 4.60                   | 217.5                  | 6.2                       | 7.84                         | 99.5                        |                               |
| Elastic plate 2 | CXF51×102   | 397,400                         | 1,589,600 0.66                | 2.64                     | 8.10                    | 8.10                   | 123.5                  | 6.2                       | 8.84                         | 67.5                        |                               |

Table 2  Estimation of dry natural frequency of the elastic walls 1 and 2 (wave impact test)

| Elastic wall | Spring type | Stiffness (single spring) (N/m) | Stiffness (four springs) (N/m) | Mass (single spring) (kg) | Mass (four springs) (kg) | T (single spring) (ms) | T (four springs) (ms) | Freq. of single spring (Hz) | Mass of impact wall (kg) | Mass of four springs + impact wall (kg) | Freq. of four springs + impact wall (Hz) |
|--------------|-------------|--------------------------------|--------------------------------|--------------------------|-------------------------|------------------------|------------------------|---------------------------|------------------------------|------------------------------------------|-------------------------------|
| Elastic wall 1 | CL51×102    | 98,493                          | 393,973                        | 0.475                    | 1.9                     | 13.8                   | 13.8                   | 72.5                      | 13.7                         | 15.6                        | 25.3                         |
| Elastic wall 2 | CL51×254    | 37,702                           | 150,807                        | 1.15                     | 4.6                     | 34.7                   | 34.7                   | 28.8                      | 13.7                         | 18.3                        | 14.4                         |
Appendix B: wave impacts on rigid and elastic plates

See appendix Figs. 21, 22, 23, 24, 25, 26.

Fig. 21 Time histories (left) and their spectra (right) of acceleration, force and deflection due to broken wave impact

Fig. 22 Pressure time histories (left) and their spectra (right) due to broken wave impact on the walls. Note the different vertical scales
Fig. 23  Time histories (left) and their spectra (right) of acceleration, force and deflection due to high aeration impact.

Fig. 24  Pressure time histories (left) and their spectra (right) due to high aeration wave impact on the walls. Note the different vertical scales.
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