LYMAN BREAK GALAXIES AT $z \sim 3$ AND BEYOND

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Abstract

We report on the status of large surveys of photometrically selected star forming galaxies at $z \sim 3$ and $z \sim 4$, with particular emphasis on both the advantages and the limitations of selecting objects using the “Lyman break” technique. Current results on the luminosity functions, luminosity densities, color distribution, star formation rates, clustering properties, and the differential evolution of the population as a function of redshift are summarized.

1 Introduction

We have heard a great deal at this meeting about the exciting results on high redshift galaxies that are beginning to be obtained at IR, sub-millimeter, and radio wavelengths. These are clearly very special times for studies of the distant universe, which are destined to unfold rapidly with the coming of ubiquitous 8m-class telescopes, new observing techniques, capabilities at previously unexplored wavelengths, and, not least, by increased awareness of what one should be looking for to successfully find galaxies at high redshift. If any message emerges when one looks at this meeting as a whole, it is that there are many complementary ways to acquire information on the distant universe, and one probably needs some synthesis of all of the observations, at all wavelengths, to appreciate the full picture of what is going on. It is important to be cognizant of the advantages and limitations of each technique, and to pay attention to what is happening throughout the field and not fall victim to “wavelength myopia”, wherein one promotes observations at some particular wavelength without accounting for what one has learned, or can learn, from other methods. In the interest of trying our best to acquire this kind of global view of the current situation, in this short summary we make an attempt to give an honest assessment of the virtues and limitations of studying high redshift galaxies by selecting them using the “Lyman break” technique.

Any high redshift object that produces UV continuum photons (whether it be produced by star formation or by AGN activity) that manage to escape the parent object without being absorbed and re-radiated in the far-IR is capable of being detected using the Lyman break technique; the guaranteed spectral feature at the ionization edge of neutral hydrogen at 912 Å in the rest frame can be discerned easily with broad-band photometry. While the idea of using this feature to find distant galaxies is not new (see [29] for an amazingly prescient discussion of the predicted appearance and possible search methods for “primeval–galaxies”),
it is only recently that it has been successfully implemented \cite{13,14,28,29}. In only a few years, the high efficiency of the technique has resulted in very rapid progress in accumulating large samples of high redshift galaxies. Large samples allow statistical insight into the nature of a particular galaxy population, and enable the same kind of detailed analyses that have recently been completed at $z \sim 1$ \cite{25,15,12,10}.

Of course, Lyman break galaxies (LBGs) are certainly not the whole story, and many of the talks at this meeting have emphasized how incomplete any assessment of the progress of galaxy formation at high redshift can be if based entirely on objects selected in the far–UV. The far–UV is easily subject to extinction by dust, and is sensitive essentially only to current (unobscured) massive star formation, and not at all to accumulated stellar populations of a galaxy (as would arguably be the case for near–IR selected samples at moderate redshift–see, e.g., Peter Eisenhardt’s contribution from this meeting.) Thus, one learns nothing about the star formation history of individual objects by studying Lyman break galaxies (although one is not obviously biased against objects that have substantial stellar mass, so long as they also have current star formation, and it is possible to make measurements sensitive to older stars once the objects are found) and it is not clear a priori to what level dust has affected what appears or does not appear in the sample.

Having stated these caveats up front, we move to the advantages of LBG selection. An obvious one is that the surface density of these high redshift galaxies is high enough that large samples can be acquired using existing CCD imagers and imaging spectrographs on 4m-10m telescopes. This has allowed the first detailed analyses of the clustering properties of galaxies at $z >> 1$ \cite{42,43,16,1}. The clustering properties are a real acid test of whether the basic ideas we have about how galaxies form are anywhere close to correct. Even a field as small as the Hubble Deep Field provides large enough samples of LBGs that many have used it to infer the entire history of galaxy and star formation \cite{28,36}, although we will argue below that this application is premature.

One of the reasons for the success of these early LBG studies is that the observed–frame optical (and therefore rest–frame UV for the high redshift objects) is still the spectral region in which the greatest sensitivity is achievable from the ground. For example, the spectroscopic surveys for LBGs that we will describe below achieve limiting sensitivities of about 200 nano-Jy at 0.7 $\mu$m. With purely photometric techniques, particularly using deep HST images, these numbers can be improved upon significantly. This large dynamic range allows the detection (in principle) of heavily obscured high redshift galaxies so long as there is a small leakage of UV photons. Put another way, the ground–based LBG samples can detect the equivalent of several $\times 10^9 L_\odot$ at $z \sim 3$, so that an object like Arp 220 could be detected even if only $\sim 1\%$ of the energy emerged in the far–UV. Now, it might not be recognized as an ultra-luminous galaxy on the basis of the far–UV observations, but there is some hope that one can disentangle the effects of extinction and reddening for the LBGs, as discussed below.

## 2 Lyman Break Galaxies at $z \sim 3$

Because LBGs have been reviewed many times recently \cite{13,25,13}, in this short summary we concentrate on new results and on general inferences, rather than on details.

At the time of this writing, we have obtained spectroscopic redshifts for more than 750 galaxies in the redshift range $2 \leq z \leq 5$, using various color–selection criteria to choose targets for spectroscopy. The most intensive work to date has been on a sample at $z \sim 3$, for which about 0.25 square degrees (in 6 high latitude fields, mostly of angular extent 9’ by 18’) has been imaged (using the facilities of many 4m–class telescopes, including Palomar, CTIO, KPNO,
Figure 1: (Left) Theoretical 2–color diagram illustrating the region from which \( z \sim 3 \) LBGs are selected. The 3 color tracks correspond to a model spectrum subjected to different amounts of reddening, spanning the range observed for most of the real galaxies: \( E(B-V)=0 \) (squares), \( E(B-V)=0.15 \) (triangles), and \( E(B-V)=0.30 \) (pentagons). The points are at intervals of \( \Delta z = 0.1 \), with redshifts between 2.7 and 3.4 enlarged. (Right) The current redshift histogram for objects selected from the indicated region of the two color diagram. The dotted curve represents the predictions based on the models shown in the left panel, accounting for known sources of incompleteness.

and WHT) deeply enough to reliably select candidates to \( R_{AB} = 25.5 \). Our spectroscopic sample has been obtained exclusively using the Keck telescopes and the Low Resolution Imaging Spectrograph [33]. The selection criteria and the current redshift histogram are shown in Figure 1.

Unlike many spectroscopic surveys, the imaging portion for target selection is absolutely crucial and is also very time–consuming: a typical field, which is imaged in the \( U_{n}GRI \) filters to a uniform depth of \( \sim 29 – 29.5 \) mag per square arc second surface brightness limits, requires more than 2 full nights of good weather and good seeing on a 4m–class telescope. Huge gains in this aspect of the project could be realized with the new generation of wide–field mosaic imagers, tiled with CCDs having good UV/blue response. As we will discuss below (§3), we have just completed a similarly large imaging survey (but smaller spectroscopic sub–sample) of galaxies at \( z \sim 4 \), using analogous selection criteria but different filter combinations for candidate selection.

2.1 Far–UV Luminosity Function

Recently we have undertaken a major effort to quantify the details of our selection criteria, and to understand the various sources of incompleteness in the ground–based color–selected samples [2, 41]. Extensive Monte Carlo simulations have been used to determine the effects on our samples of photometric errors, blending with foreground objects, and “template incompleteness” (i.e., what fraction of all galaxies at \( z \sim 3 \) would be detected given our color selection criteria). The results of these simulations provide an estimate of the effective volume surveyed as a function of apparent magnitude (essential for calculating a luminosity function) given the incompletenesses and an estimate of the true distribution of galaxy colors at a given redshift (see §2.2). The latter can then be used in conjunction with models to obtain a better estimate
Figure 2: (Left) The observed far–UV luminosity function of $z \sim 3$ Lyman break galaxies. The HDF points have been re–evaluated using newly determined effective volumes for the selection criteria outlined in [13], and no re–normalization has been necessary to make the two data sets consistent with one another. The data are plotted as a function of apparent magnitude, rather than absolute magnitude– the observed $R$ magnitude corresponds to $\sim 1700$ Å in the rest frame at $z = 3$. (Right) The inferred distribution of reddening for the $z \sim 3$ LBGs, assuming that all differences in the far–UV continua are produced by dust reddening with the reddening relation of [6].

of the effective volume surveyed for any set of LBG selection criteria, including those used in the HDF [28, 27, 13]. The resulting composite far–UV luminosity function is shown in Figure 2a.

Note that a Schechter [37] function is a very good fit to the distribution, and that the quite–steep faint end slope is obtained either by using only the ground–based data to $R = 25.5$, or by including the HDF data which reach to considerably fainter magnitudes. Most of the difference between this LF and an earlier version presented in [13], which had a much flatter faint end slope, is attributable to the greatly–improved incompleteness corrections for the ground–based samples.

Of course, the observed far–UV luminosity function has almost certainly been significantly altered by extinction. Any estimate of this effect is likely to be highly uncertain until it is possible to measure the same galaxies in the far–UV and in the far–IR (see §2.5), but the simulations mentioned above allow the construction of an internally–consistent model that is also consistent with external observations, as discussed in the next section.

2.2 Colors, Reddening, and Extinction

As shown in Figure 2b, a by–product of our large spectroscopic sample and incompleteness corrections is an estimate of the intrinsic distribution of continuum colors among the $z \sim 3$ LBGs. Here E(B-V) was simply used as a parameterization of the range of continuum colors encountered; the negative values of E(B-V) correspond to rare objects which have very strong line emission in the $G$ band, which can have as much as a 0.2 magnitude effect on the colors–none of the objects in the sample has continuum colors bluer than our assumed model spectrum against which we are measuring the implied reddening. The values of E(B-V) are dependent on our assumptions about the intrinsic LBG spectrum shape, and on our choice of reddening
law, which in this case is taken from [5]. As can be seen from the figure, most of the LBGs have implied reddening that lies in the range E(B-V) = 0 – 0.3, with a median value of 0.15, which corresponds to an implied extinction at rest-frame 1500 Å of ~ 1.7 magnitudes, or a factor of ~ 5, with a net correction in total UV luminosity across the population in the spectroscopic sample of a factor of ~ 7 (cf. [13]). Individual (rare) objects, however, have implied extinction of up to ~ 5 magnitudes. The implications are that a “typical” LBG, with an apparent magnitude corresponding to m* in the luminosity function shown in Figure 2a, and the median extinction correction, would have a star formation rate of about 65h_{50}^{-2} M_\odot yr^{-1} for Ω_M = 1; the most luminous objects in the sample, after nominal correction for extinction, would have SFR in excess of 1000 M_\odot yr^{-1} for the same cosmology.

Within our spectroscopic sample, we also see an observed color–magnitude relation, in which the more apparently luminous are also redder, on average. Thus many of the brightest objects in the sample become significantly brighter when the extinction corrections are applied, while the faintest objects tend to be bluer and to have smaller implied extinction corrections. This trend is consistent with what is observed for local galaxies (cf. [20]). The effects on the luminosity function are explored by [2].

How reliable are these extinction estimates? The UV spectral slope has been observed to be a reasonably reliable predictor of reddening for nearby UV–selected objects for which far–IR observations are also available [30], and the implied corrections based on the UV slope are quite consistent with those estimated from the Balmer emission lines for the few z ~ 3 galaxies with near–IR spectroscopy [4]. In addition, the typical implied reddening of the z ~ 3 LBGs is also consistent with that estimated for lower redshift objects for which both rest–UV and Hα observations have been compared [18, 46], after accounting for the difference in rest–UV wavelength at which the extinction is estimated. Hence, on empirical grounds, the levels of extinction implied seem reasonable. Nevertheless, the extinction in any individual case is probably quite uncertain.

These issues are discussed in much more detail in [11, 2, 13, 7, 34].

2.3 Clustering

Most of what we know at present on the clustering properties of the LBGs has already been published in a series of papers [12, 13, 16, 1, 17]. The basic result is that at z ~ 3, the relatively bright LBGs in the ground–based sample are very strongly clustered, with a co–moving correlation length r_0 that is equal to or greater than that of present–day galaxies. A comparison of the clustering strength with the expected clustering of the dark matter distribution at the same epoch indicates a significant level of bias, ranging from b ~ 2 to b ~ 6 depending on the model. Moreover, there is evidence for “clustering segregation” as a function of UV luminosity, in the sense that fainter samples of LBGs are less strongly clustered [17]. A very good match to the number density and clustering properties of LBGs and dark matter halos in simple analytic theory (e.g., [1, 33]) is found if the shape of the power spectrum is that which best accounts for present–day large scale structure, normalized by the abundance of rich clusters locally. Within the context of most of the currently popular models, the characteristic mass of LBGs in the ground–based samples is ~ 10^{12} M_\odot; the high bias of the LBGs is explained by the fact that they reside in halos that are rare at z ~ 3. A large number of theoretical papers have addressed the clustering of LBGs at high redshift (e.g., [12, 21, 3, 19, 3, 22, 31, 17, 10, 23]), essentially all finding the observations consistent with general hierarchical models if the LBGs are tracing the most massive virialized halos at z ~ 3.

In [1] we point out that the agreement of the dark matter models and the observations of real galaxies (based upon UV luminosity) in terms of abundance and clustering strength requires
that there be a monotonic relationship between UV luminosity and dark matter halo mass, with relatively small scatter. In [2] we consider the details of the implications for the relationship between mass and luminosity for LBGs given the observations and simple assumptions about the dark matter distribution.

It is likely that the largest peaks in redshift histograms such as those shown in Figure 4 are the progenitors of rich clusters of galaxies today [42, 19], and are seen prior to collapse. Thus, many of the LBGs’ descendents are likely to be found in rich environments today. It is plausible that LBGs are closely related to massive galaxies in the universe today, and that the star formation one sees at $z \sim 3$ is producing the stars presently found in the bulges of early type spirals and in ellipticals. As always, a direct connection to present–day galaxies is not straightforward, and is model–dependent. At present, there are significant differences among the various groups in the detailed model predictions (e.g. [4, 23, 40]); the differences are almost entirely due to the star formation recipes, and not to the behavior of the dark matter.

2.4 Kinematics

One of the most important pieces of information that is still missing but that is observationally tractable is the dynamical masses of LBGs. As seen above, the clustering properties already suggest association with large masses, in the context of any reasonable hierarchical structure formation model. However, as discussed in [34, 1], the actual masses of the halos that are as abundant as observed LBGs vary considerably as a function of $\Omega_M$ for a fixed power spectrum shape. A critical test of the validity of any particular structure formation scenario would be actual direct measurements of LBG masses.

The far–UV spectra of LBGs, while very interesting, are severely limited for measuring any kind of dynamical information because most of the absorption lines are produced by outflowing interstellar gas, and Lyman $\alpha$ emission is severely affected by radiative transfer effects and by external absorption. It is essential, therefore, to use spectral diagnostics that are more likely to reflect gravitationally induced motions.
Toward this end, we have attempted pilot observations targeting the nebular lines in the rest–frame optical (observed near–IR)\(^{(34)}\) in an effort to measure line widths as a crude estimate of mass. At present the observations are so difficult that the progress is painstakingly slow and the results are somewhat ambiguous from a dynamical point of view. The interpretation of line widths as indicative of masses is also fraught with uncertainty; observed line widths are essentially always an underestimate of the true circular velocity of a galaxy because most of the emission emerges from a small region that is often on the rising part of the rotation curve– see also \(^{(31)}\). Nevertheless, we plan a major program of near–IR spectroscopy of LBGs in the near future, using 8m–class telescopes. At the very least, the data can be used as a cross–check on the extinction corrections discussed above (cf. \(^{(34)}\)), and, by measuring [OIII], H\(\beta\), and [OII], estimates of chemical abundances will be feasible.

2.5 LBGs at Other Wavelengths

An obvious and important extension of the issues discussed in §2.2 above is to attempt to observe LBGs at rest–frame far–IR wavelengths, as both a cross-check on extinction estimates, and as a means of determining how much overlap there is between the LBG population discovered in the UV and the luminous IR galaxies discovered using SCUBA, discussed extensively at this meeting. The main problem is that few of the LBGs, even after extinction correction, imply star formation rates that would be detectable with SCUBA, where the detection/confusion limit is typically \(> 400 \, M_\odot \, yr^{-1}\) for the adopted cosmology. The first results on pointed sub–mm observations of LBGs will soon be available \(^{(8)}\). On the other hand, given the extinction estimates described in §2.2, a considerable fraction of the far–IR background must be produced by the equivalent of LBGs at intermediate to high redshifts (cf. \(^{(7)}\)).

In addition, we have been pursuing a program of near–IR imaging of a substantial subset of the ground–based LBG sample, and excellent near–IR data for LBGs in the HDF will soon be available (e.g., \(^{(14)}\)). These observations will improve the wavelength baseline for reddening estimates, and offer the possibility of crude assessments of the stellar populations in these early star–forming galaxies.

In general, multi–wavelength campaigns will be crucial for understanding the complete energetics and nature of LBGs.

3 Beyond \(z \sim 3\)

Within the past couple of years, inferences on the global star formation history over \(\sim 90\%\) of the age of the universe have become possible (cf., e.g., \(^{(28, 27, 36)}\)). The first estimates \(^{(28)}\) were based upon a combination of spectroscopic redshift surveys \(^{(23)}\), photometric redshifts in the HDF \(^{(11)}\), and photometric Lyman break galaxies in the HDF. Quite rightly, these first results have motivated both a great deal of excitement, and a large number of papers pointing out how the picture could very well be misleading. There are two obvious problems, one of which has been very much in the forefront at this meeting: the obvious one is that the HDF results only account for galaxies that are selected in the far–UV, and do not account for the “optically dark” component of galaxy formation (e.g. \(^{(4)}\)). We have seen above, and elsewhere, that even for galaxies selected in the UV, most of the energy is apparently absorbed and is presumably re–emitted in the far–IR. A second possible problem is that the HDF is a very small area, and it is quite possibly dangerous to reach universal conclusions based on a 2′ patch of sky, no matter how exquisite the data. While it will take considerable time to fully address the first problem,
the second one is in a regime that can be addressed directly using large-area ground-based surveys.

One of the most intriguing results from the initial forays into understanding the “star formation history of the universe” was the apparently rapid decline in the star formation density for redshifts beyond \( z \sim 3 \) [28, 27]. This result was qualitatively consistent with the decline observed for quasars at both optical [38, 24] and radio [39] wavelengths, and provided the rather satisfactory feeling that perhaps one was observing the entire epoch of galaxy formation, all the way to the beginning of the “dark ages”.

Given the very solid baseline of our \( z \sim 3 \) photometric and spectroscopic samples, we set out to assemble an equivalent and analogous sample at \( z \sim 4 \) in order to test the validity of the HDF result on the evolution of the UV luminosity density with redshift. Figure 4a illustrates the color criteria used to select the same type of objects present in the \( z \sim 3 \) sample, in order to make the cleanest possible differential comparison. The results of our preliminary spectroscopic follow-up are shown in Figure 4b. A full description of the survey and results are given in [41].

The spectroscopic results so far indicate that the population of UV–luminous galaxies at \( z \sim 4 \) is very similar to that observed at \( z \sim 3 \) – while we do not have as much basis for measuring the UV continua of the \( z \gtrsim 4 \) galaxies, the assumption that they have the same distribution of intrinsic colors as measured for the \( z \sim 3 \) population results in the predicted redshift distribution shown in Figure 4b, which is so far quite consistent with the spectroscopic redshifts. The spectroscopy, while not as extensive as for the \( z \sim 3 \) sample, allows a robust estimate of the effective volume of the survey using techniques analogous to those described above for the \( z \sim 3 \) sample, and also allows a correction for contamination by lower–redshift interlopers, which is \( \sim 20\% \) for the color criteria shown in Figure 4a. When we then compare the far–UV luminosity functions at \( z \sim 3 \) and \( z \sim 4 \), the result is shown in Figure 5a. Integrating to obtain UV luminosity density for the range of intrinsic luminosity in common for the two
samples, the result is $\rho_{\text{UV}}(z = 3)/\rho_{\text{UV}}(z = 4) = 1.1 \pm 0.3$, independent of cosmology [41]. Thus, at least at the bright end of the far–UV luminosity function, there is no significant difference in the luminosity distribution or in the normalization between $z \sim 3$ and $z \sim 4$. This result is purely empirical and depends on few assumptions.

Of obvious interest (albeit subject to many more uncertainties), given the discussions at this meeting, is how these new surveys covering volumes $\sim 200$ times larger than that of the HDF affect the inferences about the star formation history of the universe. In Figure 5b, we have plotted the new results together with earlier results, where in all cases we have integrated over the luminosity functions only down to the equivalent of $\sim 0.1L^*$ for consistency (this is an extrapolation for all but the $z \sim 3$ LBG sample). Integrating the fitted LFs to arbitrarily small luminosities would increase the high redshift points by another factor of $\sim 1.7$ relative to the lower redshift points because of the much steeper LF that applies. The extinction–corrected version was obtained by conservatively applying the median extinction correction discussed above for the $z \sim 3$ and $z \sim 4$ samples. The point we have tried to make [41] in this exercise is that the HDF may have provided misleading results, probably due to sample variance given the small volume, and at this point the much better constraints provided by the large ground–based samples are quite consistent with a universe in which the co–moving SFR density remained constant between $z \sim 4.5$ and $z \sim 1$. This could have many implications, which space constraints prevent us from discussing here; some are discussed in [41].

4 Summary

LBGs offer the distinct advantage of allowing wholesale surveys of galaxies in the distant universe, where it is straightforward to choose nearly volume limited samples of star forming galaxies in a prescribed redshift range, and where follow–up work is within reach of the current
generation of 8m–class telescopes and current instruments. The efficiency of the Lyman break technique has allowed the first robust estimates of the clustering properties of distant galaxies, and one of the most robust luminosity functions determined for any sample beyond the local universe. The clustering properties in particular have allowed for direct testing of general frameworks for galaxy formation, and the good statistics allow for detailed comparison with models of galaxy and structure formation, creating a new interface with theory that is certain to yield progress in the coming few years.

On the other hand, the LBG surveys are possibly a very incomplete picture of what is going on at high redshift. The technique almost certainly misses where most of the energy produced by star formation at high redshift is emerging (although one can attempt to account for at least some of this with extinction estimates, and it is not yet clear how many objects are being missed), and it is not yet known if the overall picture painted by LBGs for the history of star formation parallels that which will eventually emerge from studies of high redshift galaxies in the sub-mm regime. It is also not possible at present to make more than plausibility arguments connecting LBGs to present–day galaxies, and the technique is largely insensitive to star formation which has occurred prior to the epoch of observation.

A sensible and constructive combination of what we are all learning using all of the various techniques available should nicely fill in a scenario of which we now have only a thumbnail sketch. What is really amazing is how far things have come in just the last few years for understanding the nature of high redshift galaxies. With all of the capabilities becoming available in the near and not too distant future, the prospects for imminent progress are immense.

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