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Xiaowei Li1,2, Aiying Wang2 and Kwang-Ryeol Lee1

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INTRODUCTION

In recent years, amorphous carbon (a-C) films have become increasingly popular in both the scientific and engineering communities because of their impressive combination of high hardness, low friction coefficient, and chemical inertness.1–4 These properties can provide effective protection against serious mechanical or chemical damage of moving mechanical components, such as tools,5 automobile engines,6 aerospace components,7 and artificial joints.8,9 However, a-C films inevitably undergo self-consumption, which severely limits their lifetime and reliability and ultimately leads to the catastrophic failure of the coated surface.

According to recent reports, the combination of a fluid lubricant (normally a base oil with additives) and an a-C film can ensure easy slippage and long wear life of the friction interface.10–13 Thus significantly enhancing the anti-friction and anti-wear abilities of moving mechanical components over a long service lifetime. However, a majority of reports have focused on the anti-friction and anti-wear properties of inorganic nanoparticles and organic modifiers, such as zinc dialkyldithiophosphate,14,15 tallow diethanol amine,16 glycerol mono-oleate,17 and graphene,18 incorporated in the lubricant, while the interaction between the a-C film and base oil and the corresponding evolution of the interfacial structure during the friction process has not been given enough attention. Additionally, descriptions of the synergistic mechanism of the a-C and fluid films are still confined to phenomenal explanations due to the highly limited experimental in-situ characterization of the microstructure of the friction interface. Thus, an in-depth understanding of the friction mechanism at the atomic scale is required.

Furthermore, it is well known that the contact pressure strongly affects the friction properties of a-C films.2–5 In particular, in the case of an a-C composite with lubricating oil, changing the contact pressure can trigger the transition of the lubrication state between the a-C and the lubricating oil from boundary lubrication to mixed and hydrodynamic lubrication according to the classical Stribeckcurve theory.20,21 Hence, exploring the friction behavior of the a-C/lubricating oil system under different contact pressures may not only further the scientific understanding of their synergistic friction mechanism in different lubrication states, but also provide an understanding of the friction-induced structural transformation of the lubricating oil, which is indispensable for the selection and design of the base oil and the technical implementation of a-C-lubricant systems.

In the present work, we performed a reactive molecular dynamics (MD) simulation using ReaxFF22,23 to study the friction behavior of a self-mated a-C film lubricated with a base oil. The linear alpha olefin C5H10 was considered as the representative lubricant. The dependence of the friction properties of the a-C/C5H10/a-C system on the contact pressure was investigated, and the corresponding evolution of its interfacial structure was studied to unravel the fundamental mechanism of their synergistic friction behavior. The results indicated that the friction properties were highly dependent on the contact pressure, and the low friction was attributed to the synergistic effects of both interfacial passivation and C5H10 hydrodynamic lubrication.

RESULTS

Observation of morphology during the friction process Figure 1 shows the evolution of the configurations of a-C/C5H10/a-C friction systems with sliding time under contact pressures of 5, 20, and 50 GPa, respectively. As shown in Fig. 1a, when the contact pressure was 5 GPa, the system contained two sharp a-C/C5H10...
interfaces with locally mixed atoms after the loading process at 0 ps. From 0 to 1250 ps, the atomic movement at the interface was enhanced, and the sharp interface became somewhat fuzzy and eventually transformed into a steady-state friction interface. In particular, a fraction of the C₅H₁₀ molecules were still located at the central position and did not interact with the a-C structures (Fig. 1a). However, as the contact pressure was increased to 20 and 50 GPa (Fig. 1b,c), more atoms diffused across the interface, resulting in greater interaction between the a-C film and C₅H₁₀ lubricant. Hence, the initially observed sharp a-C/C₅H₁₀ interfaces could not be distinguished. A severe structural distortion occurred, suggesting the presence of high residual compressive stress at the interface. In addition, the system required a long time for the running-in process to reach the steady state because of the drastic mixing at the interface, as will be confirmed later.

Friction with respect to contact pressure for the a-C/C₅H₁₀/a-C system
The changes in contact pressure, temperature, kinetic energy (KE), and potential energy (PE) with sliding time during the friction process are plotted in Fig. 2a. First, the specified contact pressure was achieved for each case. The changes in temperature, KE, and PE, which increase as a function of contact pressure, confirmed the duration of the running-in processes. However, Fig. 2b clearly shows that the average temperature at the steady-state friction stage increased with contact pressure following a strong fluctuation (see error bars in Fig. 2b); the average temperatures were 304 ± 3 K at 5 GPa, 317 ± 6 K at 20 GPa, and 519 ± 44 K at 50 GPa. The flash temperature between both surface asperities can be estimated using the following equation:²,²⁴

\[ \Delta T = \frac{\mu W v}{8 a K_{a-C}} \]  

where \( \Delta T \) is the rise of the flash temperature at the contact between the asperities; \( \mu \), \( W \), and \( v \) are the friction coefficient, applied load, and sliding velocity, respectively; \( a \) is the radius of the real contact area, and \( K_{a-C} \) is the thermal conductivity of the a-C material. According to Eq. (1), as the sliding velocity was fixed at 10 m/s for each case, the change in the flash temperature should be dominated by the friction force, as will be discussed later. Moreover, compared to the lubricant-free system (Fig. 2b), the introduction of the C₅H₁₀ lubricant suppressed the increase in the temperature at the a-C/a-C friction interface at low contact pressures (5 and 20 GPa), while it unexpectedly increased at a high contact pressure (50 GPa), suggesting the presence of strong interactions between the a-C and the C₅H₁₀ lubricant, such as bond breaking and dissociation of C₅H₁₀ molecules. The changes in PE and KE with the contact pressure were also consistent with that of the temperature.

Figure 2c shows the evolution of the friction force and load curves with sliding time during the friction process. In each case,
Fig. 2 Friction results with respect to contact pressure for the a-C/C5H10/a-C system. a Evolution of the contact pressure, temperature, KE, and PE with sliding time for each case. b Average temperature of the free layer during the steady-state friction stage as a function of contact pressure. c Evolution of the friction force and load curves with sliding time for each case. d Average friction force and load, and e friction coefficient as a function of contact pressure. For comparison, the results for the pure a-C/a-C system without the C5H10 lubricant were also considered. Error bars were standard deviations.

The stable friction stage could be achieved after a sliding time of 1250 ps. The durations of the running-in process were further confirmed, and were consistent with the changes in the temperature and energies shown in Fig. 2a. In order to evaluate the dependence of friction on contact pressure, the friction force and load values during the last 200 ps (yellow region in Fig. 2c) were considered separately to compute the average friction force and load during the stable friction process, as illustrated in Fig. 2d. It can be noted that as the contact pressure increased from 5 to 50 GPa, the average friction force increased drastically. The rapid increase in temperature (Fig. 2b) can be explained by Eq. (1); the load also showed a trend similar to that of the friction force. However, different rates of increase of the average friction force and load with contact pressure were observed, leading to distinct changes in the friction coefficient, as shown in Fig. 2e. For comparison, the results of the pure a-C/a-C system without the C5H10 lubricant were also considered. When the contact pressure was increased from 5 to 50 GPa, the friction coefficient first decreased and then increased, following the general intrinsic shape of a Strubeck curve. The minimum friction coefficient of 0.36 was obtained for the friction system at 20 GPa. Most importantly, compared to dry conditions, the use of C5H10 as a lubricant drastically improved the friction properties of the a-C film. However, when the contact pressure was 5 GPa, the friction coefficient was reduced by 93%, while a reduction of only 19% was observed at a contact pressure of 50 GPa, suggesting that the anti-friction efficiency of C5H10 was strongly dependent on the contact state between the two sliding a-C films.

Definition of the sliding interface and intrinsic region

In order to clarify the relationship between the friction behavior and contact pressure, the friction-induced evolution of the structural properties must be analyzed. Therefore, the atomic distributions of a-C and the C5H10 lubricant and the evolution of the structural properties of the a-C/C5H10/a-C system along the z direction were first evaluated (Fig. 3). Based on the atomic distributions and structural properties along the z direction, the system can be divided into three layers in each case: the bottom intrinsic a-C, upper intrinsic a-C, and interfacial layers. At a low contact pressure of 5 or 20 GPa (Fig. 3a, b), variations in the atom number, density, residual stress, and coordination number with sliding time were mainly observed in the interfacial layer rather than the bottom and upper intrinsic a-C layers. However, at a contact pressure of 50 GPa (Fig. 3c), all the layers showed obvious transitions in their structural properties. In addition, the width of the interfacial layer decreased with increasing contact pressure, ranging from 18 Å at 5 GPa to 15 Å at 20 GPa and 12 Å at 50 GPa. In order to investigate the effect of contact pressure on the a-C/C5H10/a-C system, the interface (gray region in Fig. 3), bottom intrinsic a-C (green region in Fig. 3 with a thickness of 12 Å), and upper intrinsic a-C (blue region in Fig. 3 with a thickness of 12 Å) regions were defined separately for each case to quantify their structure and properties.

Evolution of structural properties with sliding time in each case

The dependence of the density, residual stress, and hybridized structure of the bottom intrinsic a-C layer on sliding time in the a-C/C5H10/a-C system at different contact pressures is summarized in Fig. 4. First, the evolution of the structural properties of the bottom intrinsic a-C layer was consistent with that of the upper layer (see Figure S1 of Supplementary Information). During the sliding process, almost no changes were observed in the structure and properties with sliding time at the low contact pressures of 5 and 20 GPa. However, when the contact pressure was increased to 50 GPa, the density, residual compressive stress, and sp3-C fraction as a function of sliding time increased significantly, accompanied by a drop in the sp2-C and sp-C fractions. This indicated that a high contact pressure could induce a strong sp2-to-sp3 transition.
transformation according to the $P$-$T$ phase diagram,$^{26,27}$ which could improve the wear-resistance of the a-C films; however, the coated surface could also potentially be exfoliated due to the high residual stress. The interfacial layer also showed behavior similar to that of the bottom a-C layer (see Figure S2 of Supplementary Information). Friction behavior is known to be strongly related to the structure and properties of the interfacial layer. However, since the interface is composed of a-C and C$_5$H$_{10}$ lubricant, before characterizing the structure and properties of friction interface, it must first be determined whether the a-C film binds with the C$_5$H$_{10}$ lubricant. These characterizations are a prerequisite to clarify the friction mechanism of the a-C/C$_5$H$_{10}$/a-C system at different contact pressures.

Information on the binding of C$_5$H$_{10}$ with a-C surface

Figure 5 shows the contribution of the C$_5$H$_{10}$ lubricant to the coordination number of the a-C structure and the final morphologies of the separated a-C or C$_5$H$_{10}$ lubricant in the a-C/C$_5$H$_{10}$/a-C system after the friction process. The number of intact C$_5$H$_{10}$ molecules and the bonding ratio of each C atom in the C$_5$H$_{10}$ lubricant are also evaluated in Fig. 6. The bonding ratio, $B_{ab}$, is defined as follows:

$$B_{ab} = \frac{N_{ab}^t}{N_{ab}^0}$$

where $N_{ab}^t$ is the total bond number of atom $a$ contributed by atom $b$ when the sliding time is $t$; $N_{ab}^0$ is the total bond number of
Fig. 5  Information at the a-C/C\textsubscript{5}H\textsubscript{10} friction interface. a Contribution of the C\textsubscript{5}H\textsubscript{10} lubricant to the coordination number of the a-C structure; b a-C and c C\textsubscript{5}H\textsubscript{10} lubricant structures in the a-C/C\textsubscript{5}H\textsubscript{10}/a-C system after the friction process at different contact pressures.

Fig. 6  Dissociation of C\textsubscript{5}H\textsubscript{10} lubricant. a Number of intact C\textsubscript{5}H\textsubscript{10} molecules and b bonding ratio for each of the C-C bonds of C\textsubscript{5}H\textsubscript{10} after the friction process, in which only the contribution from the lubricant is considered.
atom $a$ contributed by atom $b$ when the sliding time is 0 ps; $a$ and $b$ range from atom C2 to C3 and C4 in the C5H10 molecules. With increasing contact pressure, the contribution of C5H10 to the coordination number of the a-C structure obviously increased. As shown in Fig. 5a, when the contact pressure was 5 GPa, the C5H10 molecules made no contribution to the coordination number of the a-C structure, suggesting that there was no chemical bonding between the a-C and C5H10 molecules but instead only intermolecular interactions. Figure 3a also reveals that the C5H10 molecules were mainly distributed at the center of the interface to form a C5H10 plateau region with a width of 3 Å, which could effectively inhibit chemical interactions between the bottom and upper a-C surfaces, as shown in Fig. 5b. In addition, there was no change in the number of intact C5H10 molecules, as illustrated in Figs. 5c and 6a. This can be treated as a hydrodynamic lubrication state according to the Striebeck curve theory.21 When the contact pressure was increased to 20 GPa, the chemical interaction between the two a-C surfaces was further enhanced (Fig. 5b), and the dissociation of a few C5H10 molecules (Figs. 5c and 6a) occurred; the fragments generated by the dissociated C5H10 could bond with the a-C surface, thus contributing to the coordination number of the a-C structure (Fig. 5a). In particular, it should be noted that the increased coordination number of a-C mainly originated from the C4 and C3 atoms of C5H10. Compared with the initial values in Fig. 6b, the bonding ratios of the C3-C3 and C4-C4 bonds decreased by 1 and 3%, respectively, indicating that the scission of carbon–carbon backbones mainly occurred at the –CH2–CH= and –CH–CH2 sites. However, most of the intact C5H10 molecules remained at the center of the interface throughout the intermolecular interactions, corresponding to a mixed lubrication state.21

When the contact pressure was further increased to 50 GPa, the bottom a-C layer strongly bonded with the upper a-C layer (Fig. 5), and the temperature at the friction interface reached up to 519 ± 44 K, leading to the severe dissociation (58%) of the C5H10 molecules (Fig. 6a) and thus reaching the boundary lubrication state.21 This dissociation was mainly achieved through further scission at the –CH=CH2 and –CH2–CH=CH2 sites, with 34 and 56% reductions in their bond number, respectively, although reductions of 10% for –CH2–CH3 and 14% for –CH–CH= were also observed. Large amounts of these fragments, which are displayed in Fig. 5c (pink circles), recombined with the a-C surface. This resulted in a drastic increase in the coordination number of a-C (Fig. 5a), which also arose mainly from the C4 and C3 atoms of the C5H10 lubricant. In addition, Erdemir et al. reported that during the friction process, the dehydrogenation of linear alkenes at the friction interface was facilitated by the catalytic action of the Cu films. However, no dehydrogenation was observed in the present work (see Figure S3 of Supplementary Information) due to the absence of catalytic action at the friction surface.

Dependence of the interfacial properties on contact pressure

The density and stress at the interface, which are contributed by the a-C and C5H10 lubricant, respectively, were investigated, and the results are illustrated in Fig. 7. When the contact pressure was increased from 5 to 20 and 50 GPa, the density of a-C increased gradually due to the enhanced interaction between the bottom and upper intrinsic a-C layers, while the contribution of C5H10 to the density also increased, which contributed to the reduction of the interfacial volume at a high contact pressure. The contribution of the C5H10 to a-C structures to the residual stress also increased obviously due to the highly distorted a-C structure and C5H10 dissociation. Compared to the lubricant-free conditions, (see Figure 54 of Supplementary Information), the introduction of C5H10 lubricant into the a-C/a-C interface lowered both the density and residual stress of the friction interface; however, at the high contact pressure of 50 GPa, this difference was diminished by the dissociation and recombination of C5H10 molecules.

Evolution of the interfacial structure with contact pressure to determine the friction mechanism

In order to explore the friction mechanism, the hybridized structure of the C atoms at the interface and the mean square displacement (MSD) of the C5H10 lubricant were calculated separately (Fig. 8). The MSD is an essential tool to evaluate the mobility of the C5H10 molecules, which was estimated using the following equation:

$$\text{MSD} = \langle r^2(t) \rangle = \frac{1}{N} \sum_{i=1}^{N} [r_i(t) - r_i(0)]^2$$

where $N$ is the number of $i$ atoms in the system, $r_i(t)$ is the position of the $i$th atom at time $t$, and $r_i(0)$ is the position of the $i$th atom at $t = 0$. At a contact pressure of 5 GPa, the friction system operated via the hydrodynamic lubrication of the C5H10 molecules, leading to a lower friction coefficient compared to that of the system without lubricant. When the contact pressure was increased to 20 GPa, there was no change in the sp$^3$ fraction, but the sp$^2$ fraction increased while the sp fraction decreased. This change
resulted from the interaction between the two a-C layers and the dissociation of a few C\textsubscript{5}H\textsubscript{10} molecules, which passivated the dangling bonds at the interface. In particular, Fig. 8a shows that the MSD at a contact pressure of 20 GPa showed almost no change from that at 5 GPa, although a few C\textsubscript{5}H\textsubscript{10} molecules were dissociated. The shearing strength, \( \sigma \), is shown in Fig. 8c as a function of contact pressure during the steady-state friction process, and was calculated as follows:

\[
\mu = \frac{S}{\sigma}
\]

where \( S \) is the shearing strength (GPa) and \( \sigma \) is the Hertzian contact pressure (GPa). Hence, the synergistic effect of both the passivation of dangling bonds and C\textsubscript{5}H\textsubscript{10} hydrodynamic lubrication lowered the rate of increase of the shearing strength with increasing contact pressure (Fig. 8c), accounting for the minimum friction coefficient at 20 GPa (Fig. 2e).

For the system at a contact pressure to 50 GPa, although the dissociation of a large amount of C\textsubscript{5}H\textsubscript{10} molecules and the interaction between the two a-C layers clearly reduced both the sp\textsuperscript{2} and sp fractions (Fig. 8b), the most important factor was the weakening of the hydrodynamic lubrication from C\textsubscript{5}H\textsubscript{10} (Fig. 8d), which led to a faster increase in the shearing strength (Fig. 8c). This induced an increase of 27.8% (from 0.36 to 0.46) in the friction coefficient at 20 GPa (Fig. 2e).

The present results span the range between the organic chemistry, surface science, and applied engineering, provide a scientific understanding of the alpha olefin-lubricated a-C system, and most importantly, provide effective guidance for the selection and design of lubricants and promote the development of novel a-C-lubricant composite lubrication systems for existing and potential technological implementations.

**METHODS**

Large-scale atomic/molecular massively parallel simulator (LAMMPS) code\textsuperscript{28} was used to perform the reactive MD simulations. Figure S6 in the Supplementary Information shows the model and the related parameters used in this calculation. The a-C structure was deposited by an atom-by-atom method\textsuperscript{29} and was set as the lower and upper mating materials. The structure was composed of 6877 carbon atoms and had an sp\textsuperscript{3} fraction of 24 at.%, an sp\textsuperscript{2} fraction of 72 at.%, and a density of 2.7 g/cm\textsuperscript{3}. C\textsubscript{5}H\textsubscript{10} was selected as a representative lubricating oil. Its surface was located 3 Å away from the lower or upper a-C film, and the number of C\textsubscript{5}H\textsubscript{10} molecules was fixed at 72 for each case. A three-layer model was used for the a-C/C\textsubscript{5}H\textsubscript{10}/a-C system, including a fixed layer (blue region in Figure S6), thermostatic layer (purple region in Figure S6), and free layer (remaining region in Fig. 6) to model the evolution of the friction-induced interfacial structure. In this model, the fixed layer was used to simulate the semi-infinite system, the thermostatic layer was maintained at 300 K using the 

\[ NVE \]

ensemble with a Berendsen thermostat\textsuperscript{30} to provide a thermal reservoir for the simulated system, and the others, including the remnant a-C and C\textsubscript{5}H\textsubscript{10} lubricant, were totally free. A time step of 0.25 fs was used, and the periodic boundary condition was applied along the x- and y-directions.
The ReaxFF potential developed by Srinivasan and Tavazza et al. was selected to describe the C-C, C-H, and H-H interactions. Additional validation was also carried out by simulating the growth of a-C films using the atom-by-atom deposition approach, and the corresponding results revealed that the ReaxFF potential used in the present work could accurately describe the density, residual stress, hybridized structure, bond angles, and length distributions of the a-C structures through comparison with the widely accepted AIREBO potential (manuscript under preparation). During the friction process, a three-step process was adopted (Figure S6): (i) geometric optimization (GO) at 300 K for 2.5 ps, (ii) a loading process to achieve the specified value of contact pressure (5, 20, and 50 GPa) during 25 ps, and (iii) a sliding process with a fixed contact pressure (5–50 GPa) and sliding velocity (10 m/s) along the x-direction for 1250 ps. A high contact pressure, such as 50 GPa, was possible for the instantaneous contact of a-C asperities during the friction process. The evolution of the structural properties after the GO and loading processes can be found in Figure S7 of the Supplementary Information. After the friction process, the friction coefficient ($\mu$) was calculated as follows:

$$\mu = \frac{f}{W}$$

where the frictional force, $f$, is calculated by summing the force acting on the fixed atoms of the lower a-C model in the sliding direction, while $W$ is the load acting on the fixed atoms of the lower a-C model in the z direction. The force acting on the sliding atoms was averaged every 100 MD steps, and thus a total of 50000 individual data points for the frictional force and load were obtained during 5000000 MD steps.

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding authors, specifically Professor Kwang-Ryeol Lee of the Korea Institute of Science and Technology (email: krlee@kist.re.kr) upon reasonable request.

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AUTHOR CONTRIBUTIONS

X.L. and K.R.L. designed the calculations and performed the friction analysis. X.L., K.R., and A.W. co-wrote the manuscript. All authors participated in the discussions and manuscript preparation.

ADDITIONAL INFORMATION

Supplementary information accompanies the paper on the npj Computational Materials website (https://doi.org/10.1038/s41524-018-0111-x).

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