THE BRIGHT END OF THE ULTRAVIOLET LUMINOSITY FUNCTION AT $z \sim 8$: NEW CONSTRAINTS FROM CANDELS DATA IN GOODS-SOUTH*

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Abstract

We present new $z \sim 8$ galaxy candidates from a search over ~95 arcmin$^2$ of WFC3/IR data, tripling the previous search area for bright $z \sim 8$ galaxies. Our analysis uses newly acquired WFC3/IR imaging data from the CANDELS Multi-Cycle Treasury program over the GOODS-South field. These new data are combined with existing deep optical Advanced Camera for Surveys (ACS) imaging to search for relatively bright ($M_{\text{UV}} < -19.5$ mag) $z \sim 8$ galaxy candidates using the Lyman break technique. These new candidates are used to determine the bright end of the UV luminosity function (LF) of star-forming galaxies at $z \sim 7.2$–8.7, i.e., a cosmic age of 600 ± 80 Myr. To minimize contamination from lower redshift galaxies, we make full use of all optical ACS data and impose strict non-detection criteria based on an optical $\chi^2_{\text{opt}}$ flux measurement. In the whole search area, we identify 16 candidate $z \sim 8$ galaxies, spanning a magnitude range $H_{\text{F60,AB}} = 25.7$–27.9 mag. The new data show that the UV LF is a factor ~1.7 lower at $M_{\text{UV}} < -19.5$ mag than determined from the HUDF09 and Early Release Science (ERS) data alone. Combining this new sample with the previous candidates from the HUDF09 and ERS data allows us to perform the most accurate measurement of the $z \sim 8$ UV LF yet. Schechter function fit to the combined data result in a best-fit characteristic magnitude of $M_\ast(z = 8) = -20.04 \pm 0.46$ mag. The faint-end slope is very steep, though quite uncertain, with $\alpha = -2.06 \pm 0.32$. A combination of wide-area data with additional ultra-deep imaging will be required to significantly reduce the uncertainties on these parameters in the future.

Key words: galaxies: evolution – galaxies: high-redshift – galaxies: luminosity function, mass function

Online-only material: color figures

1. INTRODUCTION

Thanks to the unprecedented efficiency of the WFC3/IR camera on board the Hubble Space Telescope (HST), the last two years have seen remarkable progress in the exploration of galaxies within the first Gyr after the big bang. Immediately after its installation, WFC3/IR has been used to obtain the deepest NIR images ever seen as part of the HUDF09 program (PI: Illingworth, e.g., Bouwens et al. 2011b). This pushed the observational frontier of galaxies from $z \sim 6$ well into the reionization epoch, with first constraints even at $z \sim 10$ (Bouwens et al. 2011a; Oesch et al. 2012; Zheng et al. 2012).

In combination with deep ancillary optical data from the HUDF (Beckwith et al. 2006) and GOODS surveys (Giavalisco et al. 2004), the WFC3/IR data from the HUDF09 and the Early Release Science (ERS; Windhorst et al. 2011) programs allowed for the identification of more than 130 $z \sim 7$–8 galaxy candidates to date using the Lyman break galaxy (LBG) selection technique (e.g., Steidel et al. 1996; Giavalisco 2002). Since the Ly$\alpha$ absorption of the neutral intergalactic medium (IGM) shifts to $>1\mu$m at $z \gtrsim 7$, such galaxies can only be detected in the NIR data and are completely invisible in the optical. For example, $z \sim 8$ galaxies can be selected as $Y_{105}$-dropouts, based on their very red $Y_{105}-J_{125}$ colors and optical non-detections.

The first $z \sim 7$–8 galaxy candidates have already been used to measure the UV LF at these redshifts (e.g., Oesch et al. 2010c; Bouwens et al. 2010a, 2011b; McLure et al. 2010; Bunker et al. 2010; Finkelstein et al. 2010, 2011; Yan et al. 2010; Wilkins et al. 2010; Lorenzoni et al. 2011; Bradley et al. 2012). Additionally, these candidates allowed for first estimates of the physical parameters of $z > 6$ galaxies, such as their sizes (Oesch et al. 2010a), UV-continuum slopes (e.g., Bouwens et al. 2010b, 2012b; Finkelstein et al. 2010, 2012; Dunlop et al. 2012), rest-frame optical colors (Gonzalez et al. 2012), and even ages and stellar masses (e.g., Labbé et al. 2010a, 2010b; González et al. 2010, 2011; Finkelstein et al. 2010; Schaerer & de Barros 2010). Furthermore, $z \sim 8$ protocluster candidates have been identified in pure parallel WFC3/IR data (Trenti et al. 2012).

While the first deep WFC3/IR data sets provide good constraints on the faint population at $z > 7$, they only probe a limited volume and therefore result in poor constraints on the much less abundant bright galaxies around the exponential cutoff of the luminosity function (LF). At $z \sim 7$, constraints on this bright cutoff can be obtained using ground-based, wide-area data from, e.g., Subaru or the Very Large Telescope (e.g., Ouchi et al. 2009; Castellano et al. 2010a; Hickey et al. 2010). However, beyond $z \sim 7$, the galaxy population is too faint, and the sky is too bright to efficiently detect $z \gtrsim 8$ galaxies from the ground with current facilities, resulting only in upper limits at $M_{\text{UV}} < -22$ mag (Castellano et al. 2010b) in the field. Note that ground-based searches behind lensing clusters resulted in a few potential $z \gtrsim 8$ candidates (see, e.g., Laporte et al. 2011).
and in limits on the bright end of the $z \sim 9$ LF (Laporte et al. 2012).

Due to the limited volume probed at $z \sim 8$ with current deep HST data, the bright end of the $z \sim 8$ LF has thus remained quite uncertain, subject to large cosmic variance. For example, in the HUDF09-2 field data, Bouwens et al. (2011b) identified two very bright ($H_{160,AB} \sim 26$ mag) sources (see also Wilkins et al. 2011a; McLure et al. 2011), which possibly biased the previous measurement of the $z \sim 8$ UV LF toward somewhat higher values. It is therefore essential to study larger area WFC3/IR data to quantify how representative these bright candidates are of the $z \sim 8$ galaxy population.

There are a few ongoing and planned HST programs that can be used for this task. First, the Multi-Cycle Treasury program CANDELS is obtaining $Y_{105}$ data over the two GOODS fields (Giavalisco et al. 2004), resulting in a total search area for $z \sim 8$ galaxies of $\sim 300$ arcmin$^2$. Second, the pure parallel programs BORG (Trenti et al. 2011; Trenti et al. 2012) and HPIIES (Yan et al. 2011b) are designed to find rare bright $z \sim 8$ galaxy candidates, sampling a combined total of about 400 arcmin$^2$ to varying depths. Before the advent of the James Webb Space Telescope (JWST), the combination of different WFC3/IR data sets probing different scales will likely be the only way to obtain a well-sampled UV LF at $z \sim 8$ (see, e.g., Bradley et al. 2012).

In this paper, we take advantage of the full $Y_{105}$-band data obtained over the GOODS-South field as part of the CANDELS program to identify relatively bright $z \sim 8$ galaxy candidates at absolute magnitudes $M_{UV} < -19.5$ mag. We combine these with our previous $z \sim 8$ searches in the ultra-deep data from the HUDF09 as well as the ERS programs to significantly improve the sampling of the bright end of the $z \sim 8$ UV LF. With the inclusion of the CANDELS field, we essentially triple the search area for bright $z \sim 8$ galaxies in GOODS-South relative to our previous analysis, resulting in much more reliable constraints on the shape of the bright end of the LF and the luminosity at which an exponential cutoff might occur.

An accurate measurement of the bright part of the UV LF is critical for enabling meaningful comparisons with galaxy evolution models at these redshifts (e.g., Finlator et al. 2011b; Finlator et al. 2011a; Trenti et al. 2010; Lacey et al. 2011; Jaacks et al. 2012a; Dayal & Ferrara 2012; Forero-Romero et al. 2010; Salvaterra et al. 2011; Muñoz 2012). The bright end potentially contains important information on the star formation efficiencies, formation timescales, and duty cycles of galaxies within dark matter halos in the first few hundred Myr of cosmic time.

Furthermore, a well-sampled bright end of the LF helps to break degeneracies in fitting the Schechter function (Schechter 1976). Even if shallower data sets do not provide real information on the faint-end slopes, they can still decrease the uncertainties on the faint-end slope, assuming that the shape of the LF can accurately be described by a Schechter function (e.g., see discussion in Bouwens et al. 2008). Accurate estimates of the faint-end slope are critical for obtaining realistic constraints on the number of ionizing photons emitted by the ultra-faint galaxy population at these redshifts. Within the current uncertainties, it is unclear whether the faint galaxy population was luminous enough to reionize the universe alone at $z \gtrsim 7$ (e.g., Bolton & Haehnelt 2007; Oesch et al. 2009; Robertson et al. 2010; Bouwens et al. 2012a; Shull et al. 2011).

Finally, the identification of brighter $z \sim 8$ galaxy candidates is crucially important to provide candidates for spectroscopic follow-up. Spectroscopy has proven to be extremely challenging, mainly due to the intrinsic faintness of the sources, the abundance of bright night-sky lines, and the absorption of a significant fraction of the Ly$\alpha$ flux of these galaxies by the neutral IGM. Therefore, the progress in spectroscopic follow-up has been relatively slow, with only a handful of confirmed sources at $z \sim 7$ (see, e.g., Schenker et al. 2012; Pentericci et al. 2011; Ono et al. 2012). At $z \sim 8$, possibly the best chance for spectroscopic confirmation of the redshifts of candidates is with upcoming multi-object NIR spectrographs simultaneously targeting a few bright sources identified in contiguous wider area imaging.

This paper is organized as follows. After the presentation of the data in Section 2, we explain our source selection criteria and present the $z \sim 8$ candidates in Section 3. These are then used to derive constraints on the bright end of the $z \sim 8$ UV LF in Section 4. We will refer to the HST filters F435W, F606W, F775W, F814W, F850LP, F105W, F125W, F160W as $B_{435}$, $V_{606}$, $i_{775}$, $j_{814}$, $z_{850}$, $Y_{105}$, $J_{125}$, $H_{160}$, respectively. Throughout this paper, we adopt $\Omega_M = 0.3$, $\Omega_{\Lambda} = 0.7$, $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$, i.e., $h = 0.7$ (WMAP-7; Komatsu et al. 2011). Magnitudes are given in the AB system (Oke & Gunn 1983).

2. THE DATA

2.1. WFC3/IR Data over GOODS-South

In this paper, we analyze the complete public WFC3/IR data obtained over the GOODS-South field as part of the Multi-Cycle Treasury program CANDELS (PI: Faber/Ferguson; Grogin et al. 2011; Koekemoer et al. 2011) and then combine it with the deeper HUDF09 and ERS data sets. The completion of the $Y_{105}$ imaging over the CANDELS GOODS-South field (see Figure 1) has made this a unique data set for studying the bright end of the $z \sim 8$ UV LF based on a $Y_{105}$-dropout selection.

The CANDELS data are split into two parts, CANDELS-Deep and CANDELS-Wide. The first covers the central part of GOODS-South in $3 \times 5$ tiles with $\sim 3.5$ orbits in each of the filters $Y_{105}$, $J_{125}$, and $H_{160}$. These data cover $\sim 65$ arcmin$^2$, reaching down to $H_{160,AB} = 27.7$ mag. The CANDELS-Wide program consists of data in the same three WFC3/IR filters as for the CANDELS-Deep. These were fully acquired already by end of 2011 March and cover eight WFC3/IR pointings ($\sim 35$ arcmin$^2$) to $\sim 1$ orbit depth in each filter, thus reaching to $H_{160,AB} = 27.0$ mag.

In addition, we combined all the public WFC3/IR data that have been taken over these fields before 2012 March from other programs. In particular, we included the imaging data of the supernovae follow-up program of CANDELS (PI: Riess), which adds imaging over a few pointings of CANDELS-Deep. The final exposure map of all the $Y_{105}$ data included is shown in Figure 1. The central part covered by the ultra-deep HUDF09 program was omitted from our analysis of the CANDELS data in order not to replicate the bright candidates of that field given in the study of Bouwens et al. (2011b). We will use those candidates, however, when computing the total $z \sim 8$ LF in Section 4.4 (see also Section 2.2).

The WFC3/IR data were reduced using standard techniques. The pipeline-processed science frames were obtained from MAST, and were subsequently registered to the existing optical Advanced Camera for Surveys (ACS) data at a pixel scale of 0′′.06. All input images were inspected visually for satellite trails and other artifacts, such as loss of the guide star. Pixels affected by persistence from previous observations, as identified from the
Table 1

| Field               | Area (arcmin²) | $B_{435}$ | $V_{606}$ | $i_{775}$ | $z_{850}$ | $Y_{105}$ | $J_{125}$ | $H_{160}$ |
|---------------------|---------------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|
| CANDELS-Deep        | 60.2          | 27.9      | 28.1      | 27.6      | 28.5      | 27.6      | 27.9      | 28.0      |
| CANDELS-Wide        | 34.2          | 27.9      | 28.1      | 27.6      | 28.0      | 27.6      | 27.2      | 27.3      |

Notes. The limits correspond to 5σ variations in the sky flux measured in a circular aperture of 0.25 radius, i.e., no correction to total magnitude was performed. The full CANDELS-Deep is about 3.5 orbits per filter, and CANDELS-Wide one orbit per filter, leading to a difference of about 0.7 mag.

- The quoted area corresponds to the part of the image where imaging in all three WFC3/IR filters is available as well as ACS in $B_{435}$, $V_{606}$, $i_{775}$, $z_{850}$.
- $I_{814}$ is used for confirming the optical non-detection of candidates using a 2σ non-detection criterion. All candidates listed in Table 2 do have $I_{814}$ coverage.

Figure 1. WFC3/IR F105W data over the GOODS-South field that are used in this analysis. These data include all WFC3/IR data taken over this area before the end of 2011. The different fields are color coded by exposure time, with deeper areas being darker. The CANDELS-Deep covers ~65 arcmin² with ~3.5 orbits of F105W data. The outline indicates the area of the optical ACS coverage of GOODS. The CANDELS-Wide field spans 34 arcmin² and is only of 1 orbit depth. The central part covered by the HUDF is omitted from our primary analysis since it was analyzed fully by Bouwens et al. (2011b). However, we do combine the CANDELS $z \sim 8$ candidates with the ones we previously selected in the deep HUDF09 and the ERS fields (see Bouwens et al. 2011b) when constraining the $z \sim 8$ LF. For more information on the data used here, see Section 2 and Table 1.

For accurately constraining the UV LF at $z \sim 8$, it is crucial to probe the galaxy population over as large a dynamic range in luminosity as possible. In the later sections of this paper, when we compute the LF, we therefore include all the $z \sim 8$ galaxy candidates that we previously identified in Bouwens et al. (2011b). These candidates were selected from the three ultra-deep WFC3/IR pointings of the HUDF09 program (PI: Illingworth) as well as from ~40 arcmin² of WFC3/IR ERS data (see Windhorst et al. 2011). In total, these are 59 $z \sim 8$ galaxy candidates with magnitudes in the range $H_{160, AB} = 26.0–29.4$ mag over the total area of ~53 arcmin². For more information on these candidates and their selection, we refer the reader to Bouwens et al. (2011b). As we describe below, we use essentially identical selection procedures to Bouwens et al. (2011b), to ensure that we can combine the new LF results and the older ones without corrections or biases.

3. SOURCE SELECTION

3.1. Catalog Construction

Source catalogs are obtained with the SExtractor program (Bertin & Arnouts 1996), which is used to detect galaxies in the square root of the chi-square image (Szalay et al. 1999) computed from the $J_{125}$ and $H_{160}$ data. We perform matched aperture photometry on point-spread function (PSF) matched images. The colors used here are based on small elliptical apertures (1.2 Kron; these are typically nearly round and about 0.2 long axis), and total magnitudes are measured in standard 2.5 Kron apertures (typically 0.4 radius), corrected to total fluxes using the PSF encircled energy of an equivalent aperture to account for flux loss in PSF wings (typically ~0.15–0.2 mag).
The input rms maps were scaled to properly represent the flux variations in 0′′.25 radius apertures, which were used to establish the detection significance of sources. The rms scaling was based on the detected variation in 1000 random apertures per WFC3/IR frame on empty sky regions after 3σ clipping. This procedure ensures that the weight maps correctly reproduce the actual noise in the images. Subsequently, only sources with signal-to-noise ratios (S/N) larger than 4.5 in \(H_{160}\) are considered (as measured in 0′′.25 radius circular apertures).

### 3.2. \(z \sim 8\) Color–Color Selection

Thanks to the strong IGM absorption in the rest-frame UV, \(z \sim 8\) galaxies can be selected in broadband color–color diagrams. At \(z \gtrsim 7\), the Ly\(\alpha\) absorption shifts into the \(Y_{105}\) band, rendering galaxies red in \(Y_{105} - J_{125}\), while the \(J_{125} - H_{160}\) color of star-forming galaxies is expected to be blue. The selection is discussed here, but the most challenging aspect is to minimize the contamination from low-redshift objects, as we discuss in Sections 3.3 and 3.5.

We adopt the same color selection criteria for our \(z \sim 8\) selection as in Bouwens et al. (2011b), which will allow us to directly combine the previous data with the candidates identified here. Specifically, the selection criteria are

\[
\begin{align*}
(Y_{105} - J_{125}) &> 0.45 \\
(J_{125} - H_{160}) &< 0.5.
\end{align*}
\]

Additionally, we require sources to be detected at 4.5σ in both \(J_{125}\) and \(H_{160}\), and at 5σ in at least one of these two bands. Finally, we require candidates to be detected at <2σ in all the optical ACS data, and also require candidates to meet an optical \(\chi^2_{\text{opt}}\) flux constraint. Both tests play a major role in removing potential low-redshift contaminants (see Section 3.3). Finally, we checked the \(I_{814}\) images of all sources and additionally required galaxies to be 2σ non-detections also in \(I_{814}\). This last test helped to eliminate three sources which are likely contaminants.

The above color criterion is illustrated in Figure 2, where we show the expected colors of star-forming galaxies using 100 Myr old stellar population models from Bruzual & Charlot (2003). These are additionally reddened with small amounts of dust using the dust law of Calzetti et al. (2000). The IGM absorption is modeled using Madau (1995). Our color criteria start to select such galaxies \(z \gtrsim 7.2\). The selection then peaks around \(z \sim 8\) and falls off due to the combination of an increased distance modulus and flux reduction as the Lyman break approaches the filter bandpass limit at higher redshift (see Figure 4 in Bouwens et al. 2011b). The mean redshift of our sample is \((z) = 7.9\), with 80% of galaxies expected to lie at \(z = 7.2−8.7\) (see Figure 5).

In Figure 2, we also show the colors of lower redshift galaxies using Coleman et al. (1980) templates, as well as cool dwarf stars (Burgasser et al. 2004), which could potentially contaminate our selection. Despite some overlap in this color–color diagram with \(z > 7\) galaxy candidates, we expect the contamination of such sources to be small due to our strict optical non-detection criteria (see next section and discussion in Section 3.5). This is illustrated in the right panel of Figure 2, where we show the expected \(\chi^2_{\text{opt}}\) values (defined in the next section) as a function of \(J_{125} - H_{160}\) color for different galaxy types. Most intermediate-redshift galaxies are expected to be well detected in the optical data, resulting in \(\log \chi^2_{\text{opt}} > 1\). Only completely quiescent galaxies would lie close to our selection box. The figure also includes the measured values of the CANDELS \(z \sim 8\) candidates (see Section 3.4). The sizes of the symbols...
represent their $Y_{105} - J_{125}$ colors. Our selection includes only three galaxies with relatively blue $Y_{105} - J_{125}$ colors with positive $\chi^2_{\text{opt}}$ values between one and three. These would have the highest probability of being lower redshift contaminants.

As a further step against interlopers, we check each candidate for bright detections in the Spitzer IRAC 3.6 μm images available over GOODS-South (e.g., Dickinson et al. 2003). This has proven to be a very effective test for removing dusty contaminants in very high-$z$ galaxy searches in which only limited information on the UV-continuum slope of galaxies is available (see, e.g., Oesch et al. 2012). Young, star-forming galaxies at $z > 7$ with small amounts of dust reddening are expected to show colors $H_{160} - [3.6] \lesssim 1$ (see also Gonzalez et al. 2012). Any galaxy with substantially redder colors is therefore likely to be a lower redshift contaminant.

Indeed, in constructing our sample, we identified one bright ($H_{160,AB} = 24.3$ mag) edge-on spiral galaxy (at 03:32:14.72,−27:46:21.6), which satisfied our WFC3/IR color and optical non-detection criteria. However, this source is extremely luminous in the IRAC data, with $H_{160} - [3.6] = 2.3$, and is also clearly detected even in the [8.0] band, as expected for a heavily dust-obscured source at intermediate redshift. We therefore removed it from the potential $z \sim 8$ galaxy sample. All other sources where we could measure IRAC fluxes (i.e., which were not completely blended with bright foreground sources) were consistent with the limit $H_{160} - [3.6] < 1$.

### 3.3. Optical Non-detections Using $\chi^2_{\text{opt}}$

As indicated above, one of the main challenges in selecting robust LBGs is to remove intermediate-redshift contaminants. Given the high efficiency of WFC3/IR, the available ancillary optical data in most of the WFC3/IR search fields are not appreciably deeper than the new IR data (see, e.g., Table 1). Thus, at the faintest magnitudes, even a 2σ non-detection criterion, which is the standard that is used in LBG selections, is not sufficiently effective to eliminate faint, dusty interlopers.

In previous papers (e.g., Bouwens et al. 2011b), we have developed an efficient method for eliminating low-$z$ contaminants by making full use of all the information in the optical data. In particular, for each galaxy, we compute an optical pseudo $\chi^2_{\text{opt}}$ value from its aperture flux measurements of all optical bands as

$$\chi^2_{\text{opt}} = \sum_i \text{SGN}(f_i(f_i/\sigma_i)^2,$$

where $i$ runs over $B_{435}$, $V_{606}$, $i_{775}$, and $z_{850}$, and SGN is the sign function, i.e., $\text{SGN}(x) = -1$ if $x < 0$ and $\text{SGN}(x) = 1$ if $x > 0$.

For real high-redshift candidates, the $\chi^2_{\text{opt}}$ distribution is expected to be centered around zero, while for contaminants the distribution is skewed toward positive values. Therefore, by removing galaxies with values above some limiting value $\chi^2_{\text{lim}}$, it is possible to significantly reduce the number of contaminants in LBG samples.

The adopted limiting value is derived from two different sets of simulations. The results are shown in Figure 3 where we compare the observed $\chi^2_{\text{opt}}$ to those from our simulations to establish the limiting value.

The first simulation is based on the HUDF09 data, where the available optical and WFC3/IR data are $\sim 1.5–2$ mag deeper than in the fields studied here. At a given $H_{160,AB}$ magnitude, this allows us to check what fraction of sources would contaminate our sample, if they were observed at the shallower depth of our data. We use a Monte Carlo simulation in which we apply Gaussian scatter to the fluxes of the HUDF09 sources scaled to the depth of the CANDELS-Deep and Wide data. We then compute the $\chi^2_{\text{opt}}$ values for all sources that did not satisfy our selection criteria in the ultra-deep data, but which would have been selected and would have passed the optical 2σ non-detection criterion. The distribution of $\chi^2_{\text{opt}}$ values of these contaminants is shown in Figure 3 as an orange histogram.

For the second simulation, we use brighter sources directly from the CANDELS fields which are dimmed to fainter magnitudes. Thanks to the wider area relative to the HUDF, this allows us to probe rarer interloper populations and so is complementary to the HUDF09 simulation. We only use sources with $H_{160,AB} = 24–25.5$ mag as input, in order to keep the population as close as possible to the magnitude range of our candidates. Again, for each source satisfying the color selection after applying Gaussian noise, we compute the value of $\chi^2_{\text{opt}}$. This is shown in Figure 3 as a dark red line.

As is apparent from the figure, the $\chi^2_{\text{opt}}$ distribution of interlopers is significantly skewed to positive values relative to regions of empty sky as expected for real $z \sim 8$ sources. After some consideration, we adopted a limit $\chi^2_{\text{opt}} < 3.0$ as being a good discriminator between $z \sim 8$ galaxies and contaminating sources. By using this limit of $\chi^2_{\text{opt}} < 3.0$, it is possible to reduce the contamination rate by a factor of $\sim 3–4$ relative to a simple 2σ non-detection criterion in all optical bands (see also Appendix D of Bouwens et al. 2011b). At the same time, we lose just $\sim 15\%$ of real $z \sim 8$ sources. We correct for this small loss in the estimate of the selection volume.

The expected residual number of contaminants, after applying the $\chi^2_{\text{opt}}$ limit, is $1.0 \pm 0.5$ sources in the CANDELS-Deep data, and $<0.1$ in the whole CANDELS-Wide field. Based on these
simulations, we thus estimate a contamination rate that is only \(\sim 10\%\) for our full sample. Note that, given our ignorance of the LF of interlopers and cosmic variance in the input galaxy sample for our photometric scatter simulations, it is conceivable that the true contamination rate is as high as 15\%–20\%.

3.4. The Bright CANDELS \(z \sim 8\) Candidates

Applying all the selection criteria outlined above, we identify 16 new \(z \sim 8\) galaxy candidates in the full CANDELS dataset; 14 in CANDELS-Deep, and two in CANDELS-Wide. Their properties are listed in Table 2, and individual images for all candidates in the WFC3/IR bands as well as a stack of all the optical ACS images are shown in Figures 4 and 12.

The candidates span a range of \(H_{160,AB} = 25.7–27.9\) mag. Our brightest source is thus among the most luminous \(z \sim 8\) candidates known to date. It is \(\sim 0.5\) mag brighter than all other CANDELS sources, and would be an ideal target for future spectroscopic follow-up.

Such a luminous object raises an interesting possibility. If we assume that the UV luminosity of galaxies scales with halo mass, then such a bright \(z \sim 8\) galaxy is expected to lie in an overdense environment with \(\sim 5\) fainter galaxies within a diameter of 60 arcsec (see, e.g., Trenti et al. 2012; Muñoz & Loeb 2008). Unfortunately, the source lies close to an extremely bright star, reducing the detectability of fainter companions significantly. Nevertheless, we would have expected to see a few fainter companions around this source. None are found, however. Unfortunately, the proximity to the bright star also makes it impossible to obtain IRAC flux measurements from the current Spitzer data to strengthen or refute the high-redshift solution of this source. Deeper optical or IR data (or a spectrum) may thus be needed to fully resolve the origin of our brightest candidate.

As can be seen from Figure 4, all \(z \sim 8\) candidates are quite compact, but appear to be resolved. Their size distribution is consistent with the expectations based on extrapolation from lower redshift samples out to \(z \sim 7\) (e.g., Oesch et al. 2010a). Only two sources show signs for elongated morphologies, indicative of disk structures. However, the surface brightness limit in the CANDELS data work against clearly detecting such structures, compared to e.g., the HUDF (see, e.g., Oesch et al. 2010a).

At the same time that our paper was first submitted, the CANDELS team also submitted a catalog of \(z \sim 8\) galaxy candidates in Yan et al. (2011a) from the CANDELS-DEEP area. A significant fraction of their sources (45\%) does not satisfy our strict optical non-detection criteria, with a few showing clear optical detections, thus disqualifying them for being at \(z > 7\). All of their candidates are discussed in detail in the Appendix, where we also tabulate our measurements for these sources.

3.5. Sample Contamination

Thanks to our restrictive non-detection criteria, we do not expect high contamination levels in our samples. Nevertheless, some residual contamination cannot be excluded. In particular, we briefly discuss several potential sources of contamination below.

1. Low-mass stars. As shown in Figure 2, ultra-cool dwarf stars can exhibit similarly red colors as the potential high-redshift galaxies. Given the resolution of WFC3/IR, however, we can directly check whether any of our sources are unresolved. This test can reliably be done only for the brighter candidates, with well sampled profiles. Indeed, we did remove one source from our candidate list with colors consistent with a T-dwarf which appeared to be point like (at 03:32:25.33, \(-27:48:54.2\)). Furthermore, the contamination by dwarf stars is also expected to be low based on their observed surface densities in high galactic latitude fields (e.g., Ryan et al. 2005, 2011).
2. **Photometric scatter.** The most probable source of contamination in our sample comes from photometric scatter of sources with intrinsic colors similar to high-z galaxies. This has already been quantified in Section 3.3. To summarize, after applying the \( \chi^2_{\text{opt}} < 3.0 \) cut, we only expect \( 1.0 \pm 0.5 \) contaminants due to photometric scatter in the CANDELS-Deep data set, and <0.1 in the CANDELS-Wide. Therefore, with our strict optical non-detection criteria, we limit the amount of contamination to \( \sim 10\% \).

Similar conclusions are reached by investigating the \( \chi^2_{\text{opt}} \) distribution. As can be seen in the right panel of Figure 2, our sample includes five galaxies with \( \chi^2_{\text{opt}} > 0.9 \). Statistically, we expect only 22% of galaxies to be above such values, i.e., 3.5 sources. This would suggest again a contamination of \( \sim 10\% \) in the full sample.

3. **Spurious sources.** All of our candidates are very well detected both in \( H_{160,AB} \) as well as in \( J_{125} \). The chance for spurious >4.5\( \sigma \) detections at the same location in both bands is negligibly small. Additionally, most of the sources are detected at lower significance in \( Y_{105} \), further reducing the chance of a spurious source. Spurious sources are clearly not a concern for our sample.

4. **Transients.** Since the optical data has been taken a few years prior to the new WFC3/IR data, it is possible that supernovae that went off in the meantime are selected as infrared detections without any optical counterparts. However, since one of the specific goals of CANDELS-Deep is to search for such supernovae, the \( H_{160} \) and \( J_{125} \) data acquisition has been distributed over several different epochs. So far, eight epochs have been acquired. This allows us to check directly whether any of our sources could potentially be a supernova. We therefore group the individual epochs in four bins sorted by exposure date and check whether all of our sources are still visible in these sub-splits (in a \( H_{160} + J_{125} \) image). All CANDELS-Deep candidates are indeed detected in these images. For the two candidates in CANDELS-Wide data, this test is not as decisive, as the \( J_{125} \) and \( H_{160} \) data acquisition has
only been split over two epochs. Furthermore, we stress again that these sources do not appear to be point like. We therefore conclude that supernovae have not contaminated our sample.

4. THE z ∼ 8 LF

We will now use the z ∼ 8 galaxy candidates identified in the previous section to derive constraints on the z ∼ 8 UV LF, and combine these new results with previous estimates to give the best-available LF at z ∼ 8.

4.1. Selection Functions and Redshift Distributions

To compute the LF, we first have to estimate the completeness, C(m), and redshift selection functions, S(z,m). Following Oesch et al. (2007, 2009), this is done by inserting artificial galaxies with varying magnitudes, profiles, and sizes in the observational data and rerunning the source detection with the exact same setup as for the original catalogs. This is done for each of the fields individually.

Our simulations are based on using real galaxies at lower redshift and scaling them to higher redshifts using well-established evolutionary and cosmological relationships, i.e., using the “cloning” methodology of Bouwens et al. (2003). In particular, we “clone” z ∼ 4 LBGs from the GOODS and HUDF fields to higher redshifts. The images of these z ∼ 4 sources are scaled to the desired input magnitude, and are stretched to account for the difference in angular diameter distance, as well as a size scaling of (1 + z)−1 as observed for the LBG population across z ∼ 3–7 (see, e.g., Ferguson et al. 2004; Bouwens et al. 2004; Oesch et al. 2010a). This procedure ensures that the distribution of morphologies and profiles of the simulated galaxies is as close to reality as possible and thus specifically accounts for SB-dimming effects, which result in losing larger, resolved galaxies from the samples. The cloned galaxies are then inserted into the observed images with galaxy colors as expected for star-forming galaxies between z = 6 and z = 9.5 (see also Figure 2). The adopted colors are based on a UV-continuum slope distribution of β = −2.5 ± 0.4 motivated by recent determinations of the UV-continuum slopes as a function of UV luminosity at z > 6 (see, e.g., Bouwens et al. 2009, 2010b; Stanway et al. 2005; Finkelstein et al. 2010, 2012; Wilkins et al. 2011b; Dunlop et al. 2012). As a cross-check, we have also tested that our simulation pipeline returns essentially equivalent results when adopting theoretical Sérsic galaxy profiles (see also Bouwens et al. 2011b; Oesch et al. 2012).

From the simulation output, we compute the completeness as a function of the observed H_{160,AB} magnitude for each field, taking into account the scatter and bias between input and output magnitudes. Additionally, we compute the selection probabilities as a function of redshift and magnitude by measuring the fraction of sources that meet our selection criteria. The selection function and the corresponding redshift distribution function are shown in Figure 5. As can be seen, galaxies in our sample are selected from z ∼ 7.2 to 8.7, with a mean redshift of ⟨z⟩ = 7.9.

We note that the expected redshift distribution has wings that extend to lower and higher redshifts, and we expect ∼4 objects to lie at z < 7.5. The shape of the selection function has to be taken into account when modeling the observed number densities, particularly if there is strong evolution over the redshift range z = 7–8. The redshift selection function is available from the authors upon request.

4.2. Expected Surface Density of z ∼ 8 galaxies

It is very instructive to compare the expected surface density of z ∼ 8 galaxy candidates using the previous best-fit LF of Bouwens et al. (2011b) with the observed number of candidates in the CANDELS data. The expected number of sources in a given magnitude bin m_i can be estimated for any given LF, φ(M), through

\[ N_i^{\text{exp}} = \int_{\Delta m} dm \int dz \frac{dV}{dz} S(m, z)C(m)\phi(M[m, z]). \]

In Figure 6, we show the histograms of the observed surface density of sources in the different fields and compare them to the expectation from different LFs. As can be seen, the expected source density peaks around 0.2 arcmin^{-2} mag^{-1} in the CANDELS-Deep field. The peak occurs at H_{160,AB} = 27 mag, which is ∼0.5 mag brighter than the formal 5σ limiting magnitude. This is mainly due to the reduction in the selection volume due to photometric scatter, which causes us to lose sources from the color–color selection window near the limit. Note that we
also expect a small fraction of sources even below the formal magnitude limit due to scatter in the photometric offsets between the aperture fluxes (from which the selection $S_u/N$ is computed) and the total fluxes (used in Figure 6), as well as from variations in the depth of the $H_{160}$-band data.

From the integration of the surface density expected from the best-fit LF of Bouwens et al. (2011b), we find that we would expect to detect 22 $z \sim 8$ galaxies in the CANDELS-Deep field, and five in CANDELS-Wide. This is a factor $1.7 \times$ higher than the 16 candidates we find, indicating that the bright end of the previous UV LF might have been estimated somewhat high due to the presence of a few very bright sources ($H_{160, AB} \sim 26$ mag) in a possible overdensity in the HUDF09-2 field (see also the discussion in Bouwens et al. 2011b). The difference seems large but it is driven by just a few bright sources and so is consistent with small number statistics.

4.3. New Constraints on the $z \sim 8$ LF

We now use the new CANDELS $z \sim 8$ candidates to derive direct constraints on the UV LF, by computing the stepwise LF in bins of absolute magnitudes. This is done using an approximation of the effective selection volume as a function of observed magnitude $V_{\text{eff}}(m) = \int_0^\infty dz dV/dz S(z, m) C(m).$ The LF is then given by $\phi(M) dM = N_{\text{obs}} / V_{\text{eff}}(m).$ This is shown in Figure 7, where we evaluated the LF in bins of 0.6 mag, as a compromise between luminosity resolution and $S/N$. The error bars include a 30% (40%) contribution from cosmic variance for CANDELS-Deep (CANDELS-Wide), which we estimate using the cosmic variance calculator of Trenti & Stiavelli (2008); see also Robertson (2010).

As expected from the low observed surface density of $z \sim 8$ candidates in the CANDELS data, the stepwise LF is significantly below the previous best-fit Schechter function from Bouwens et al. (2011b). At $M_{UV} = -21$ to $-19$ mag, we find values $\sim 2-4 \times$ lower than from the HUDF09 and ERS data. As noted above, these are still consistent with the current and previous uncertainties. Nevertheless, the CANDELS data indicate that the characteristic magnitude might be fainter than previously determined. We quantify this more precisely in the next section.

Finally, we also compare our LF determination with the bright end constraints from the BORG survey (Bradley et al. 2012), which are based on 33 $z \sim 8$ galaxy candidates with $J_{125} < 27.4$ mag identified over an effective search area of 274 arcmin$^2$ (i.e., about $1.7 \times$ the area of the HUDF09+ERS+CANDELS). As Figure 7 shows, the BORG $z \sim 8$ LF also supports the indication that the HUDF09 field is somewhat overdense in bright $z \sim 8$ galaxies. The LF from BORG is a factor $\sim 2-3 \times$ lower than the stepwise LF determination from the HUDF09+ERS alone, and it is in excellent agreement with our new, total stepwise LF.

4.4. Combination with Deeper Data: The $z \sim 8$ Schechter Function

Due to the small dynamic range in luminosities, the bright candidates identified in the CANDELS data are not sufficient to provide a good estimate of the overall shape of the UV LF alone. We therefore combine our new sources with all of the $z \sim 8$ candidates from ultra-deep field measurements from Bouwens et al. (2011b) to update the Schechter function parameters of the $z \sim 8$ UV LF. By doing so, we can now generate an LF using a total of 75 $z \sim 8$ candidates identified over 148 arcmin$^2$, spanning $H_{160, AB} \sim 25.7$–29.5 mag. The stepwise determination of this combined sample is listed in Table 3 and shown in Figure 7.
The Schechter function parameters are derived by maximizing the Poissonian likelihood for observing \(N_{\text{obs}}\) sources in a given magnitude bin when \(N_{\text{exp}}\) are expected to be seen based on a given UV LF. We thus maximize \(L = \prod_j \prod_i P(N_{\text{obs}}^{ji}, N_{\text{exp}}^{ji})\), where \(j\) runs over all fields, and \(i\) runs over the different magnitude bins, and \(P\) is the Poissonian probability. The expected number of sources are computed according to the equation from Section 4.2. For the HUDF09 and ERS fields, we adopt the selection volumes estimated in Bouwens et al. (2011b). The combination of these older results with our new sample is appropriate since both are based on essentially identical simulations and inputs, i.e., our estimates of the CANDELS selection functions match those of Bouwens et al. (2011b).

The best-fit parameters are determined by a grid search over Schechter function parameters, maximizing the combined likelihood \(L\). The best-fit solutions are: \(\log(\Phi_\ast) = -3.30^{0.38}_{-0.46}\), \(M_\ast = -20.04^{+0.44}_{-0.48}\) mag, and \(\alpha = -2.06^{+0.35}_{-0.28}\). These parameters are consistent with the previous determination from the HUDF09 and ERS data alone (Bouwens et al. 2011b), and, despite a somewhat larger sample size, the uncertainties are not markedly reduced. However, the best-fit characteristic magnitude is fainter by \(\sim 0.1\) mag, which is due to the lower number of detected sources in the CANDELS data than the number that was expected from the previous LF determination.

Note also that the faint-end slope is steeper than \(\alpha = -2\), as discussed by Bouwens et al. (2012a). This leads to a formally divergent luminosity density. However, galaxies are not expected to be formed below a given luminosity due to inefficient cooling in low-mass halos and feedback effects (at masses that correspond to \(M_{11}\) about \(-10\) to \(-11\)), and so the luminosity density converges. Nevertheless, such steep slopes have important consequences for reionization by galaxies as pointed out previously (e.g., Bouwens et al. 2012a).

At the bright end (\(M_{11} < -20\)), our best-fit LF is a factor \(\sim 1.7\) lower than the previous determination of Bouwens et al. (2011b), who used only the HUDF09 and ERS data. However, Bouwens et al. (2011b) only find such a higher surface density for \(z \sim 8\) sources in the HUDF, while the number counts over the ERS are actually in excellent agreement with what we find over the CANDELS field. Therefore, all the wide-area data over the CDF-South GOODS yield approximately the same surface density of bright \(z \sim 8\) galaxies. Nevertheless, we stress again that both the stepwise and the best-fit LFs are consistent with the previous determination of Bouwens et al. (2011b), given the still small number of sources at the bright end.

In Table 4, we compare our new best-fit LF parameters with other, previous determinations from the literature (Bouwens et al. 2010a, 2011b; McLure et al. 2010; Lorenzoni et al. 2011; Trenti et al. 2011; Bradley et al. 2012). Within the current measurement uncertainties, these are all consistent with each other.

Overall, it is very reassuring that different groups arrive at similar results given the variety of approaches. Nonetheless, it is worthwhile remarking that with the exception of Bouwens et al. (2011b) and Bradley et al. (2012), the previous determinations were not based on large enough data sets such that all three Schechter function parameters could reliably be fit simultaneously as done here.

4.5. Evolution of the LBG Population at \(4 < z < 8\)

A key diagnostic in studying the build-up of galaxies is how the UV LF evolves with redshift, as this is directly related to the distribution of star formation rates (SFRs) in galaxies. In Figure 8, we show the error contours of the Schechter function parameters from current HST data of LBGs at different redshifts. The error contours for the \(4 \lesssim z \lesssim 7\) samples are taken from Bouwens et al. (2007, 2011b), while the \(z \sim 8\) contours correspond to the likelihood contours of our Schechter function fit at \(z \sim 8\) from all available data, i.e., including the CANDELS GOODS-South candidates (Section 4.4).

The Schechter function parameters show significant evolution from \(z \sim 7\) to \(z \sim 8\). In particular, the combination of the characteristic magnitude and number density is evolving at \(\sim 99\%\) significance. On the other hand, the combined constraint on the faint-end slope and the characteristic luminosity is still quite weak (right panel of Figure 8). Unfortunately, the faint-end slope is largely unconstrained. This is very unfortunate since it is one of the most important parameters for assessing the contribution of galaxies to reionization. The flux density of ionizing photons from galaxies is extremely sensitive to the faint-end slope \(\alpha\), especially when the slope is as steep as \(\alpha \sim -2\) (e.g., Bouwens et al. 2012a). A continued effort in the future will thus be to constrain this parameter better with deeper WFC3/IR imaging.

Overall, the UV LF constraints of LBGs are consistent with pure luminosity evolution from \(z \sim 8\) to \(z \sim 4\). This can easily be achieved by growing SFRs of individual galaxies (e.g., Stark et al. 2009; Finlator et al. 2011b; Papovich et al. 2011; Smit et al. 2012; Jaacks et al. 2012b). The best-fit evolution of \(M_\ast\) of LBGs at \(z > 3\) follows:

\[ M_\ast(z) = -20.98(\pm 0.04) + 0.31(\pm 0.03) \times (z - 3.8) \]

This is shown in Figure 9, where we plot the evolution of the characteristic cutoff luminosity of the UV LF with redshift. After correcting for dust extinction, the SFR of an L* galaxy thus grows by almost an order of magnitude from \(z \sim 8\) to \(z \sim 4\).
The value of $M_*$ as measured at $z \approx 8$ is essentially equal to the cut-off UV luminosity at $z \approx 1.5-2$, just shortly after the peak of the cosmic SFR density. However, the SFR of an $L_*$ galaxy is nevertheless larger at $z \approx 1.5-2$ than at $z \approx 8$ due to the larger dust obscuration (e.g., Reddy et al. 2010; Bouwens et al. 2012b; Smit et al. 2012).

In Figure 10, we show how the (dust-corrected) SFR density evolves across redshift when including our new UV LF parameters and the updated dust corrections for LBGs at $z > 4$ from Bouwens et al. (2012b). The SFR density includes all galaxies down to a fixed flux limit of $M_{UV} = -17.7$, which is the current detection limit at $z \approx 8$. Our updated SFR density at $z \approx 8$ is only marginally lower than the previous measurement of Bouwens et al. (2011b) by 0.07 dex, i.e., we derive $\log \rho_{SFR} = -2.32 \pm 0.12 M_\odot \text{yr}^{-1} \text{Mpc}^{-3}$. This is based on the luminosity density of $\log \rho_L = 25.58 \pm 0.12 \text{erg s}^{-1} \text{Hz}^{-1} \text{Mpc}^{-3}$ and the conversion of the UV luminosity to SFR by Madau et al. (1998).

As can be appreciated from Figure 10, the SFR density grows rather dramatically by more than an order of magnitude from $z \approx 10$ (Oesch et al. 2012; Bouwens et al. 2011a) to our new determination at $z \approx 8$. After that, it evolves very steadily, growing by another ~1.5 dex from $z \approx 8$ to its peak at $z \approx 2.5$. The rapid growth at $z > 8$ is very intriguing but still very uncertain. The change at $z > 8$ may well be refined with planned deep F140W data over the HUDF.

5. SUMMARY

In this paper, we analyzed newly acquired WFC3/IR data over the GOODS-South field as part of the CANDELS MCT program to derive new constraints on the UV LF at $z \sim 8$, about 600–750 Myr after the Big Bang. Unlike at $z \sim 7$, where the brightest galaxies can be detected from the ground, $z \sim 8$ galaxies are too faint to be detected with any reliability with
current ground-based instrumentation. Therefore, the bright end of the $z \sim 8$ UV LF has remained relatively uncertain since it has been based on WFC3/IR data that were only available over a relatively small area, the very deep HUDF09 fields.

Galaxies at $z \sim 8$ are identified using the Lyman break technique on optical and near-IR data, and particularly the $Y_{105}$ band, leading to them being identified as “$Y_{105}$-dropouts.” The CANDELS fields over GOODS are extremely valuable data sets, as they combine imaging in three WFC3/IR filters, and in addition also have relatively deep ancillary optical data from HST ACS, which allow for such $z \sim 8$ galaxy selections. Our selection criteria are outlined in Section 3 and were chosen to match prior measurements in the ultra-deep HUDF09 data set. They select galaxies at $z \sim 7.2$–8.7 with a mean redshift ($z$) = 7.9.

In the full search area of $\sim 95$ arcmin$^2$, we identified 16 new $z \sim 8$ galaxy candidates with $H_{160,AB}$ magnitudes in the range 25.7–27.9 mag. These are presented in Figure 4 and Table 2. The $z \sim 8$ candidates were selected using very strict optical non-detection requirements, including a measurement of the optical $\chi_{\text{opt}}^2$ flux (see Section 3.3). This allows us to reduce the expected contamination due to photometric scatter and due to the limited depth of the optical data relative to the WFC3/IR imaging by a factor $\sim 3$–4. Our sample is thus expected to show relatively low levels of contamination, which we estimate to be $\sim 10\%$.

Interestingly, the observed surface density of $z \sim 8$ galaxies is $\sim 1.7 \times$ lower than expected from the previous best-fit $z \sim 8$ LF of Bouwens et al. (2011b), though the difference is consistent with the statistics from the small numbers of objects at bright magnitudes in the previous search. This previous estimate was based on the much smaller area data from the HUDF09 fields and the ERS data. By using the current CANDELS imaging, we triple the search volume for luminous galaxies at $M_{\text{UV}} < -19.5$ mag in the CDFS relative to our previous analysis, thus reducing the potential biases introduced by cosmic variance.

We combine our new CANDELS $z \sim 8$ candidates with our previous candidates from the much deeper HUDF09 and ERS fields to derive the best possible measurement of the UV LF. The best-fit $z \sim 8$ UV LF we derive in this way is consistent with the previous estimates from Bouwens et al. (2011b). However, the best-fit characteristic magnitude is fainter by $\sim 0.1$ mag ($M_{*} = -20.04 \pm 0.46$ mag). Despite the larger area probed here, the uncertainties on the Schechter function parameters are still significant. Nonetheless, these new results are contributing to our growing understanding of the evolution of key parameters in the LF from $z \sim 8$ to $z \sim 2$, and of the SFR density from $z \sim 10$ to $z \sim 2$.

A combination of future wide-area data, as well as deeper WFC3/IR imaging, will be necessary to further improve on these very important measurements that are crucial for refining our estimates of the role of UV photons from galaxies in the reionization of the universe before the advent of JWST.

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Facilities: HST (ACS/WFC3), Spitzer (IRAC)

APPENDIX

COMPARISON TO YAN ET AL. (2011b) CATALOG

At the same time that our paper was submitted, the CANDELS team submitted a very similar $z \sim 8$ galaxy search in Yan et al. (2011a). That paper is still under review and is likely to change further (Yan 2012, private communication). Below, we explicitly compare only to version 2 of their manuscript, which is the version that is publicly available on arXiv at the time of publication of our paper. The Yan et al. (2011a) analysis is not based on the complete GOODS-South CANDELS data set (they include 83% of the final data in $J_{125}$ and $H_{160}$ that were used in our analysis). Additionally, they only analyze the CANDELS-DEEP area (see Figure 1), and they only include 8-epoch mosaics of the $I_{814}$ image.

For comparison, our analysis is based on all the $I_{814}$ data available over the GOODS-South, which was obtained from a large number of additional programs (ERS, HUDF09, CANDELS SNe follow-up, 3D-HST, and the UVUDF). As noted in the main text, these data reach almost 1 mag deeper than the GOODS $i_{775}$ data, and are thus an extremely valuable addition to check for low-redshift contaminants. Additionally, we include a deeper reduction of the other ACS images including additional data taken after the original GOODS-South program, which reaches to $\sim 0.1$–0.2 mag deeper compared to the GOODS v2.0 data.

The revised v2 manuscript of Yan et al. (2011a) lists a total of 16 different $Y_{105}$-dropout candidates, which are split in two different samples, one using AUTO fluxes and one using isophotal (ISO) fluxes for computing colors. Out of these 16 sources, only 5 are in common with our sample. Yan et al. (2011a) use somewhat more restrictive color criteria than what we adopt here ($Y_{105} - J_{125} > 0.8$, instead of $>0.45$ as used in our analysis). Therefore, we would have expected to select all their sources in our catalog.

Upon inspection of their sources, we found the main reason that the Yan et al. (2011a) $z \sim 8$ candidates did not appear in our catalog is due to a large fraction of their candidates showing non-negligible flux in the optical ACS data such that they do not satisfy our strict optical non-detection criterion $\chi_{\text{opt}}^2 < 3$. This applied to 7 out of their 16 candidates. Yan et al. (2011a) state that they perform a two-stage source selection, which includes a visual inspection of their candidates using all the data. However, some sources show detections in our reduction of the optical ACS data, disqualifying these sources from being at $z > 7$. Furthermore, two sources are detected in the HUDF ACS data. Since these data were used in their visual inspection, we do not have an explanation for this discrepancy.

In addition to these seven sources with optical flux, their list also includes two sources which are not significantly detected ($<5\sigma$) in our $H_{160}$ image, as well as two sources for which we measure $Y-J$ colors that are too blue, although they are consistent with being $Y$-dropouts at the $1\sigma$ level.

A detailed summary of all our measurements and analysis of the Yan et al. (2011a) sample is provided in Table 5, and an example of optical stamps for two sources is shown in Figure 11.

Given this disagreement, it is encouraging, however, that Yan et al. (2011a) do, in fact, include five out of the six sources with
Figure 11. Two examples of clear optical detections in the Yan et al. (2011a) sample (ISO_160 and ISO_164). Both sources show a detection in $I_{814}$, and show hints of flux in other bands as well. The increased depth of the $I_{814}$ data is clearly an extremely valuable addition for removing contaminants. Note that both these sources were flagged as contaminants in our catalog based on our $\chi^2_{opt} < 3$ criterion, showing the power of using this measurement.

Figure 12. Same as Figure 4, but without the slight smoothing of the stamps. Smoothing more accurately displays a larger dynamic range in the images and thus shows the reliability of the sources more fairly.
S/N(H160) > 7 from our sample which satisfy their stricter color criterion of Y−J > 0.8 (see Table 2). This indicates that our selection is quite robust, at least for higher-significance sources, if optically detected interlopers are excluded properly, as done for our catalog.

We stress again that making full use of all the information in the optical data is extremely important for a reliable LBG selection. We did this by including a limit in the optical χ²opt measurement and by analyzing all the available optical ACS data that were taken over this field. Furthermore, we note that as long as the contamination fractions and detection efficiencies are modeled self-consistently, using shallower data should not result in a different estimate of the final LF (within the errors).

REFERENCES

Arnouts, S., Schiminovich, D., Ilbert, O., et al. 2005, ApJ, 619, L43
Beckwith, S. V. W., Stiavelli, M., Koekemoer, A. M., et al. 2006, AJ, 132, 1729
Bertin, E., & Arnouts, S. 1996, A&AS, 117, 393
Bolton, J. S., & Haehnelt, M. G. 2007, MNRAS, 382, 325
Bouwens, R., Broadhurst, T., & Illingworth, G. 2003, ApJ, 593, 640
Bouwens, R. J., Illingworth, G. D., Blakeslee, J. P., Broadhurst, T. J., & Franx, M. 2004, ApJ, 611, L1
Bouwens, R. J., Illingworth, G. D., Franx, M., & Ford, H. 2007, ApJ, 670, 928
Bouwens, R. J., Illingworth, G. D., Franx, M., & Ford, H. 2008, ApJ, 686, 230
Bouwens, R. J., Illingworth, G. D., Franx, M., et al. 2009, ApJ, 705, 936
Bouwens, R. J., Illingworth, G. D., Labbe, I., et al. 2011a, ApJ, 735, 92
Bouwens, R. J., Illingworth, G. D., Oesch, P. A., et al. 2012a, ApJ, 750, L8
Bouwens, R. J., Illingworth, G. D., Oesch, P. A., et al. 2012b, ApJ, 754, 83
Bradley, L. D., TRENTI, M., OESCH, P. A., et al. 2012, arXiv:1204.3641
Bruzual, G., & Charlot, S. 2003, MNRAS, 344, 1000
Bunker, A., Wilkins, S., Ellis, R., et al. 2010, MNRAS, 409, 855
Burgasser, A. J., McElwain, M. W., Kirkpatrick, J. D., et al. 2004, AJ, 127, 2856
Calzetti, D., Armus, L., Bohlin, R. C., et al. 2000, ApJ, 533, 682
Castellano, M., Fontana, A., Boutsia, K., et al. 2010a, A&A, 511, A20
Castellano, M., Fontana, A., Paris, D., et al. 2010b, A&A, 524, A28
Cimann, G. D., Wu C., & Weedman, D. W. 1980, ApJS, 43, 395
Cucciati, O., Tresse, L., Illingworth, G. D., et al. 2012, A&A, 539, A33
Dayal, P., & Ferrara, A. 2012, MNRAS, 421, 2568

Dickinson, M., Giavalisco, M., & GOODS Team, 2003, In The Mass of Galaxies at Low and High Redshift, ed. R. Bender & A. Renzini (Berlin: Springer), 224
Dunlop, J. S., McLure, R. J., Robertson, B. E., et al. 2012, MNRAS, 420, 901
Ferguson, H. C., Dickinson, M., Giavalisco, M., et al. 2004, ApJ, 600, L107
Finkelstein, S. L., Papovich, C., Giavalisco, M., et al. 2010, ApJ, 719, 1250
Finkelstein, S. L., Papovich, C., Salmon, B., et al. 2012, ApJ, 756, 164
Finlator, K., Davè, R., & Ozel, F. 2011a, ApJ, 743, 169
Finlator, K., Oppenheimer, B. D., & Davé, R. 2011b, MNRAS, 410, 1703
Froocoer-Romero, J. E., Yepes, G., Gottlöber, S., et al. 2010, MNRAS, 403, L31
Giavalisco, M. 2002, ARA&A, 40, 579
Giavalisco, M., Dickinson, M., Ferguson, H. C., et al. 2004, ApJ, 600, L103
González, V., Bouwens, R., Labbé, I., et al. 2012, ApJ, 755, 148
González, V., Labbé, I., Bouwens, R. J., et al. 2010, ApJ, 713, 115
González, V., Labbé, I., Bouwens, R. J., et al. 2011, ApJ, 735, L34
Groot, N. A., Koevski, D., Faber, S. M., et al. 2011, ApJS, 197, 35
Hickey, S., Bunker, A., Jarvis, M. J., Chiu, K., & Bonfield, D. 2010, MNRAS, 404, 212
Jaacks, J., Choi, J.-H., Nagamine, K., Thompson, R., & Varghese, S. 2012a, MNRAS, 420, 1606
Jaacks, J., Nagamine, K., & Choi, J.-H. 2012b, MNRAS, accepted (arXiv:1204.8466)
Koekemoer, A. M., Faber, S. M., Ferguson, H. C., et al. 2011, ApJS, 197, 36
Komatsu, E., Smith, M. K., Dunkley, J., et al. 2011, ApJS, 192, 18
Labbé, I., González, V., Bouwens, R. J., et al. 2010a, ApJ, 716, L103
Lacey, C. G., Baugh, M. C., Frenk, C. S., & Benson, A. J. 2011, MNRAS, 412, 1828
Laporte, N., Pelló, R., Hayes, M., et al. 2012, A&A, 542, L31
Laporte, N., Pelló, R., Schaerer, D. et al. 2011, A&A, 531, A74
Lorenzoni, S., Bunker, A. J., Wilkins, S. M., et al. 2011, MNRAS, 414, 1455
Madau, P. 1995, ApJ, 441, 18
Madau, P., Pozzetti, L., & Dickinson, M. 1998, ApJ, 498, 106
McLure, R. J., Dunlop, J. S., Cirasolo, M., et al. 2010, MNRAS, 403, 960
McLure, R. J., Dunlop, J. S., de Ravel, L., et al. 2011, MNRAS, 418, 2074
Muñoz, J. A. 2012, J. Cosmol. Astropart. Phys., 04, 015
Puñaz, J. A., & Loeb, A. 2008, MNRAS, 385, 2175
Oesch, P. A., Bouwens, R. J., Carollo, C. M., et al. 2010a, ApJ, 709, L21
Oesch, P. A., Bouwens, R. J., Carollo, C. M., et al. 2010b, ApJ, 725, L150
Oesch, P. A., Bouwens, R. J., Illingworth, G. D., et al. 2010c, ApJ, 709, L16
Oesch, P. A., Bouwens, R. J., Illingworth, G. D., et al. 2012, ApJ, 745, 110
Oesch, P. A., Carollo, C. M., Stiavelli, M., et al. 2009, ApJ, 690, 1350
Oesch, P. A., Stiavelli, Carollo, C. M., et al. 2007, ApJ, 671, 1212
Oke, J. B., & Gunn, J. E. 1983, ApJ, 266, 713
Ono, Y., Ouchi, M., Mobasher, B., et al. 2012, ApJ, 744, 83
Ouchi, M., Mobasher, B., Shimasaku, K., et al. 2009, ApJ, 706, 1136
Papovich, C., Finkelstein, S. L., Ferguson, H. C., Lotz, J. M., & Giavalisco, M. 2011, MNRAS, 412, 1123
Pentericci, L., Fontana, A., Vanzella, E., et al. 2011, ApJ, 743, 132
Reddy, N. A., Erb, D. K., Pettini, M., Steidel, C. C., & Shapley, A. E. 2010, ApJ, 712, 1070
Reddy, N. A., & Steidel, C. C. 2009, ApJ, 692, 778
Robertson, B. E. 2010, ApJ, 713, 1266
Robertson, B. E., Ellis, R. S., Dunlop, J. S., McLure, R. J., & Stark, D. P. 2010, Nature, 468, 49
Ryan, R. E., Jr., Hathi, N. P., Cohen, S. H., & Windhorst, R. A. 2005, ApJ, 631, L159
Ryan, R. E., Thorman, P. A., Yan, H., et al. 2011, ApJ, 739, 83
Salvaterra, R., Ferrara, A., & Dayal, P. 2011, MNRAS, 414, 847
Schaerer, D., & de Barros, S. 2010, A&A, 515, A73
Schechter, P. 1976, ApJ, 203, 297
Schenker, M. A., Stark, D. P., Ellis, R. S., et al. 2012, ApJ, 744, 179
Shull, M., Harness, A., Trenti, M., & Smith, B. 2011, ApJ, accepted (arXiv:1108.3334)
Smit, R., Bouwens, R. J., Franx, M., et al. 2012, ApJ, 756, 14
Stanway, E. R., McMahon, R. G., & Bunker, A. J. 2005, MNRAS, 359, 1184
Stark, D. P., Ellis, R. S., Bunker, A., et al. 2009, ApJ, 697, 1493
Steidel, C. C., Giavalisco, M., Pettini, M., Dickinson, M., & Adelberger, K. L. 1996, ApJ, 462, L17
Szalay, A. S., Connolly, A. J., & Szokoly, G. P. 1999, AJ, 117, 68
Trenti, M., Bradley, L. D., Stiavelli, M., et al. 2011, ApJ, 727, L39
Trenti, M., Bradley, L. D., Stiavelli, M., et al. 2012, ApJ, 746, 55
Trenti, M., & Stiavelli, M. 2008, ApJ, 676, 767
Trenti, M., Stiavelli, M., Bouwens, R. J., et al. 2010, ApJ, 714, L202
Wilkins, S. M., Bunker, A. J., Ellis, R. S., et al. 2010, MNRAS, 403, 938
Wilkins, S. M., Bunker, A. J., Lorenzoni, S., & Caruana, J. 2011a, MNRAS, 411, 23
Wilkins, S. M., Bunker, A. J., Stanway, E., Lorenzoni, S., & Caruana, J. 2011b, MNRAS, 417, 717
Windhorst, R. A., Cohen, S. H., Hathi, N. P., et al. 2011, ApJS, 193, 27
Yan, H., Finkelstein, S. L., Huang, K.-H., et al. 2011a, arXiv:1112.6406
Yan, H., Windhorst, R. A., Hathi, N. P., et al. 2010, Res. Astron. Astrophys., 10, 867
Yan, H., Yan, L., Zamojski, M. A., et al. 2011b, ApJ, 728, L22
Zheng, W., Postman, M., Zitrin, A., et al. 2012, Nature, 489, 406