MHD Slip Flow and Heat Transfer of Cu-Fe$_3$O$_4$/Ethylene Glycol-Based Hybrid Nanofluid over a Stretching Surface

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Abstract: In this article, the impact of hybrid nanoparticles on different physical quantities in a Cu-Fe$_3$O$_4$/ethylene glycol-based hybrid nanofluid is associated with a steady and fully developed natural convective flow over a stretching surface. The investigation's significant results are that the ferrous oxide/ethylene glycol-based hybrid nanofluid enlarged with partial slip parameter undermines the tangential velocity and liquid suction. It causes a minute radial velocity along with temperature distribution through a stretching surface. The analysis is presented in dimensionless form. The transformed equations are solved numerically using Fourth order R-K Fehlberg with shooting technique. It is a phenomenon found in a mixture of mobile particles that exhibit specific responses to temperature strength. The particle moves to the hot cold region in thermal diffusion; then it is called 'positive'; otherwise, it is called 'negative'. The consequences of this investigation are of significance with evaluating the impact of some essential design parameters on heat transfer and, therefore, in the enhancement of industrial processes.

Keywords: MHD; hybrid nanofluid; thermal radiation; chemical reaction; porous medium; shooting technique.

List of symbols:

| $u, v, w$ | $x, y, z$ velocity components (ms$^{-1}$) |
|---|---|
| $Re$ | Reynolds number |
| $T$ | The temperature of the fluid (K) |
| $T_x$ | Free stream temperature (K) |
| $T_w$ | Surface temperature (K) |
| $C_p$ | Specific heat constant pressure (J kg$^{-1}K^{-1}$) |
| $Pr$ | Prandtl number |
| $Ec$ | Eckert number |
| $B_0$ | Strength of magnetic field (T) |
| $q_r$ | Radiative heat flux (W.m$^{-2}$) |
| $Sc$ | Schmidt number |
| $M$ | Magnetic number |
| $C$ | The concentration of the species |

| $\phi_1, \phi_2$ | Nanoparticles volume fraction |
| $\rho$ | Fluid density (kg m$^{-3}$) |
| $\mu$ | Dynamic viscosity (kg m$^{-1}$s$^{-1}$) |
| $\nu$ | Kinematic viscosity (m$^2$s$^{-1}$) |
| $\sigma$ | Electric conductivity (Sm$^{-1}$) |
| $\kappa$ | Thermal conductivity (W m$^{-1}$K$^{-1}$) |
| $\xi$ | Rate of chemical reaction |
| $\Omega$ | Angular velocity |

$^{0}$ Base fluid
$^{nf}$ Nanofluid
$^{hnf}$ Hybrid nanofluid
1. Introduction

A wide range of research has recently been carried out on a new generation of nanofluids; such liquids are known as hybrid nanofluids. Hybrid nanofluids have enhanced thermal properties because of a combined impact made by the presence of various substances. This recent development, created by the inclusion of various nanoparticles in the working liquid, has become a common topic among researchers because of its many potential benefits over nanofluids. Many researchers have currently attempted the hybrid nanofluids or mixed, combined or mixed with two different types of nanoparticles. Essentially, the idea of using hybrid nanofluids is to increase thermophysical properties in addition, particularly the charge of the heat switch compared to single-particle nanofluids or standard nanofluids. Since hybrid nanofluids are the latest form of nonliquids, the assessment process for deciding their performance remains within the improvement phase. By using hybrid nanofluids, high performance can be expected in terms of heat transfer. Recently, extended classes of nanofluids have attracted researchers’ attention, forming a variety of nanoparticles such as metal particles (Ag, Au, Al, Cu, and Fe), non-metal particles (CuO, TiO$_2$, Al$_2$O$_3$, SiC, and Fe$_3$O$_4$); carbon nanotubes. Distinct unitary hybrid nanofluids, nanofluids display uncommon rheological performance and thermo-physicochemical properties. Most significantly, hybrid nanofluids can result in suitable and acceptable nanofluid stability. Those strong nanofluids have proper functions together with better thermal conductivity at lesser nanoparticle concentrations. It also has helpful biomedical production applications, welding, lubrication, power systems, spacecraft, aeronautical device making, and vehicle thermal coordination. The application of nanomaterials for heat transfer enhancement is being widely experimented with by various researchers worldwide. At better heat flux circumstances, conventional liquids cannot achieve the desired heat transfer rate due to their inherently low thermal conductivities. The modern technique for enhancing heat transfer by utilizing ultrafine solid particles in liquids has been widely experimented with over the last decade. Choi [1] proposed the term nanofluid to describe the suspension with high effective thermal conductivity in the base fluid through the diffusion of some nanosized solid particles.

Several experimental types of research have been carried out to investigate the heat transfer properties of hybrid nanoliquids. Suresh et al. [2-3] experimentally studied convective heat transfer turbulent flow of Cu-$\text{Al}_2\text{O}_3$/H$_2$O hybrid nanofluids via a circular tube under constant heat flux conditions. They observed that the average increase in Nusselt number for Cu-$\text{Al}_2\text{O}_3$/H$_2$O hybrid nanofluid was 8.02% compared to base fluid (H$_2$O). The experimental findings show that 0.1% of Cu-$\text{Al}_2\text{O}_3$/H$_2$O hybrid nanofluids have a friction factor significantly slightly higher than 0.1% of $\text{Al}_2\text{O}_3$/H$_2$O nanofluids. Atul Bhattad [4] experimentally investigated the exergy–energy characteristics of Al$_2$O$_3$ –MgO hybrid nanofluid, mixed with various base liquids in a plate heat exchanger of the counter flow category in volume with a mixing ratio of 0.01, a volume of flows. The hybrid nanofluid operating temperatures are 30 -
45°C and 2.0 -4.0 L/min, respectively. Their results showed that the performance index of heat exchanger growths by 11.1% for alumina-MgO (4:1) deionized H2O hybrid nanofluid. Ataei et al. [5] investigated the heat transfer and flow properties of a hybrid nanofluid (i.e., TiO2-Al2O3/ H2O) in a mini channel heat source. They described an increase in the convective heat transfer coefficient compared to pure water by 16.97%. The use of TiO2-Al2O3–H2O hybrid nanofluid instead of H2O reduced the wall temperature up to 5°C. Phanindra et al. [6] employed laminar heat transfer of a hybrid nanofluid Al2O3–Cu/Oil under a boundary condition of continuous heat flux through a concentric tubular exchanger with a volume concentration of 0.1%. They claimed that the hybrid nanofluid presented the most development of 12.06% in the heat transfer rate at a Reynolds number of 1820 compared to base fluid (pure Oil).

It can be seen that many researchers considered a different type of hybrid nanofluids to enhance thermal conductivity. Taza Gul et al. [7] investigated the thin-film flow of the H2O-based carbon nanotubes hybrid nanofluid. Ogunseye et al. [8] examined the viscous dissipation on an Eyring-Powell Cu-Al2O3/ethylene glycol nanofluid. Liaquat Ali Lund et al. [9] addressed the MHD flow of Cu–Fe3O4/H2O hybrid nanofluid. Hayat et al. [10] analyzed on the enhancement in heat conduction by employing Ag–CuO/H2O hybrid nanofluid. Mabood et al. [11] examined the Al2O3–Cu–/H2O hybrid nanofluid flow with melting heat transfer. Aamir Ali et al. [12] studied enhancing the thermal conductivity of hybrid nanofluid with Cu–Al2O3 nanoparticles. Zhixiong Li et al. [13] studied the hydrothermal behavior and the irreversibility of Ag-MgO/H2O hybrid nanofluid flow numerically through a sinusoidal hairpin heat-exchanger. Benkhedda et al. [14] carried out the combined convective heat transfer improvement using TiO2-Ag/ H2O hybrid nanofluid in a heated horizontal ring. Hassan et al. [15] studied the homotopic analysis to elucidate the effects of an inclined magnetic field and nonlinear thermal radiation on the melting energy transmission in stagnation point flow over a stretchable sheet. Rahman et al. [16] carried out the transient MHD flow of water-based nanofluids along a straight-up solid porous wall against the horizon with the time-dependent plate velocity, temperature, and concentration. They computed the results by fixing the Prandtl number at $Pr=0.71$ (air), which is not valid for water-based nanofluids. Their results are unrealistic. Takabi and Shokouhmand [17] studied the rate of molecular species diffusion effects on MHD gravity-driven convective flow of nanofluids over an oscillating vertical plate Laplace transform method. The results were presented for inappropriate magnitudes of Schmidt number.

The involvement of magnetohydrodynamic in boundary layer flow controls the flow of the fluid by connecting fluid flow with magnetic fields. There has been much research on magnetohydrodynamic flow in the past. Magnetohydrodynamic flow is described as studying the magnetic properties and behavior of electro-conductive fluids through a magnetic field [18-29].

Based on the literature review, the present paper is to provide an impact of viscous dissipation and thermal radiation on MHD hybrid nanofluid flow with the combined effects of melting heat transfer and partial slip has not been previously studied. The analysis is presented in dimensionless form. The transformed equations are solved numerically using Fourth order R-K Fehlberg with shooting technique. Using these computations of the tangential velocity, temperature, and concentration have been plotted graphically to analyze the effects of various physical parameters in the problem. Also, the physical interpretation examined the skin-friction coefficient. The rate at which heat affects the boundary layer's vicinity is also presented and
discussed in detail. New insights from this study will be useful in several areas of nanotechnology and nanostructure.

2. Materials and Methods

2.1. Mathematical formulation.

The problem’s physical model is considered the 3D steady MHD hybrid nanofluid flow over a stretching sheet. The fluid occupies $z > 0$, and the flow is induced due to the stretching of the sheet $z = 0$ along the $x$– and $y$–axis with the velocities $U_w = a x$ and $V_w = b y$ and, where $a > 0$, $b > 0$ are the stretching rates of the sheet (see Fig.1). The surface temperature and concentration are subjected to partial slip, and so is $T_w + l \frac{\partial T}{\partial z}$, $C_w + l \frac{\partial C}{\partial z}$ where $l$ is the thermal slip factor. The nanofluid is turning about the vertical $z$– axis with the angular velocity $\Omega$ of the fluid is stable.

| Physical properties | Pure water | Ethylene glycol (C$_2$H$_6$O$_2$) | Cu(copper) | Fe$_3$O$_4$(Iron oxide) |
|---------------------|------------|----------------------------------|------------|-----------------------|
| $\rho$ (kg/m$^3$)   | 997.1      | 1115                             | 8933       | 5180                  |
| $c_p$ (J/kgK)       | 4179       | 2430                             | 385        | 670                   |
| $\kappa$ (W/mK)     | 0.613      | 0.253                            | 401        | 9.7                   |
| $Pr$                | 6.8377     | 24.4                             | -          | -                     |

A static magnetic field $B_0$ is imposed normal (perpendicular) to the stretching surface. It also assumed that the magnetic Reynolds number is not significant in the free convection flow; consequently, the induced magnetic field is ignored. The surface temperature is adjusted by the convection process through the hotted fluid with temperature $T_f$, $T_\infty$ and represents the ambient temperature. Ethylene glycol is the base fluid. The silver and copper are used as nanoparticles.

| Particle shapes | Spherical | Cylinder | Platelet |
|-----------------|-----------|----------|----------|
| $n$             | 3         | 4.8      | 5.7      |

Figure 1. Geometry of the problem.
The equations of the flow problem are defined as:

\[
\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0, \quad (1)
\]

\[
u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} - 2\Omega v = \frac{1}{\rho_{\text{inf}}} \left[ \mu_{\text{inf}} \frac{\partial^2 u}{\partial z^2} - \sigma B_0^2 u - \frac{\mu_{\text{inf}}}{k} u \right], \quad (2)
\]

\[
u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} + 2\Omega u = \frac{1}{\rho_{\text{inf}}} \left[ \mu_{\text{inf}} \frac{\partial^2 v}{\partial z^2} - \sigma B_0^2 v - \frac{\mu_{\text{inf}}}{k} v \right], \quad (3)
\]

\[
u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} = \alpha_{\text{inf}} \frac{\partial^2 T}{\partial z^2} - \frac{1}{(\rho C_p)_{\text{inf}}} \left( \frac{\partial q}{\partial z} \right) + \frac{\mu_{\text{inf}}}{(\rho C_p)_{\text{inf}}} \left[ \left( \frac{\partial u}{\partial z} \right)^2 + \left( \frac{\partial v}{\partial z} \right)^2 \right], \quad (4)
\]

\[
u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} + w \frac{\partial C}{\partial z} = \beta_{\text{inf}} \frac{\partial^2 C}{\partial z^2} - \xi (C - C_\infty)^n. \quad (5)
\]

The endpoint conditions are

\[
u = U_w + k v_j \frac{\partial u}{\partial z}, \quad v = V_w + k v_j \frac{\partial v}{\partial z}, \quad w = 0, \quad T = T_w + l \frac{\partial T}{\partial z}, \quad C = C_w + l \frac{\partial C}{\partial z} \quad \text{at} \quad z = 0,
\]

\[
u \to 0, \quad v \to 0, \quad T \to T_w, \quad C \to C_\infty \quad \text{as} \quad z \to \infty. \quad (6)
\]

Radiative heat flux is approximated using Rosseland approximation in the following from

\[
q_r = -\frac{4\sigma}{3\chi} \frac{\partial T^4}{\partial z} \quad (7)
\]

Where, \(\sigma\) and \(\chi\) are the “Stephan-Boltzman constant” and the mean absorption coefficient

\[
T^4 = T_w^4 + 4T_w^3 (T - T_w) + 6T_w^2 (T - T_w)^2 + ......................... \quad (8)
\]

Neglecting higher-order terms in the above equation (8) beyond the first degree in \((T - T_w)\), we get

\[
T^4 \approx 4TT_w^3 - 3T_w^4 \quad (9)
\]

Thus, substituting Eq. (9) into Eq. (7) we get

\[
\frac{\partial q_r}{\partial z} = -\frac{16\sigma T_w^3}{3\chi} \frac{\partial T}{\partial z} \quad (10)
\]

Therefore equation (4) reduced to

\[
u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} = \alpha_{\text{inf}} \frac{\partial^2 T}{\partial z^2} + \frac{16\sigma T_w^3}{3\chi} \frac{\partial^2 T}{\partial z^2} + \frac{\mu_{\text{inf}}}{(\rho C_p)_{\text{inf}}} \left[ \left( \frac{\partial u}{\partial z} \right)^2 + \left( \frac{\partial v}{\partial z} \right)^2 \right] \quad (11)
\]

By introducing the following dimensionless quantities:

\[
v = ayq' (\eta), \quad u = axp' (\eta), \quad w = -\sqrt{ay} (p (\eta) + q (\eta)) \quad \text{at} \quad \eta = z \sqrt{\frac{a}{v_f}}, \quad \theta (\eta) = \frac{T - T_w}{T_w - T_w}, \quad (12)
\]

\[
R(\eta) = \frac{C - C_\infty}{C_w - C_\infty}.
\]

2.1.1. Thermophysical properties of nanofluid and hybrid nanofluid.

Thermophysical properties are defined as follows:

The density, heat capacitance, effective thermal conductivity, and viscosity of the nanofluid are given by:

\[
\rho_{\text{nf}} = (1 - \phi) \rho_f + \phi \rho_s;
\]
\[
\left( \rho C_p \right)_{nf} = (1-\phi)\left( \rho C_p \right)_f + \phi \left( \rho C_p \right)_s
\]
\[
\kappa_{nf} = \frac{\kappa_f + (n-1)\kappa_f - (n-1)\phi \left( \kappa_f - \kappa_s \right)}{\kappa_f + (n-1)\kappa_f + \phi \left( \kappa_f - \kappa_s \right)} \text{, where } n = 3/\psi
\]
\[
\mu_{nf} = \frac{\mu_f}{\left(1-\phi\right)^{2.5}}
\]

Thermophysical properties of Cu-Fe:O₄ -Ethylene glycol nanofluid are defined as follows:
The density, heat capacitance, effective thermal conductivity, and viscosity gave by:
\[
\rho_{nf} = \left[ (1-\phi) \left( \rho_f + \phi \rho_{nf} \right) \right] + \phi \rho_{nf};
\]
\[
\left( \rho C_p \right)_{nf} = \left[ (1-\phi) \left( \rho C_p \right)_f + \phi \left( \rho C_p \right)_s \right] + \phi \left( \rho C_p \right)_s;
\]
\[
\kappa_{nf} = \frac{\kappa_{nf} + (n-1)\kappa_f - (n-1)\phi \left( \kappa_f - \kappa_s \right)}{\kappa_{nf} + (n-1)\kappa_f + \phi \left( \kappa_f - \kappa_s \right)};
\]
\[
\mu_{nf} = \frac{\mu_f}{\left(1-\phi\right)^{2.5}}
\]
The above governing equations are reduced to
\[
p''(\eta) - (1-\phi)^{2.5} (1-\phi_2)^{2.5} \left[ (1-\phi_2) \left( 1-\phi_2 \right) + \phi_2 \left( \frac{\rho s_1}{\rho f} \right) \right] + \phi_2 \left( \frac{\rho s_2}{\rho_f} \right) \]
\[
\times \left[ \left( p'(\eta) \right)^2 - p''(\eta) \left( p(\eta) + q(\eta) \right) - 2\gamma \delta q'(\eta) \right] - (1-\phi)^{2.5} (1-\phi_2)^{2.5} M^2 p'(\eta) - Kp'(\eta) = 0
\]
\[
q''(\eta) - (1-\phi)^{2.5} (1-\phi_2)^{2.5} \left[ (1-\phi_2) \left( 1-\phi_2 \right) + \phi_2 \left( \frac{\rho s_1}{\rho f} \right) \right] + \phi_2 \left( \frac{\rho s_2}{\rho_f} \right) \]
\[
\times \left[ \left( q'(\eta) \right)^2 - q''(\eta) \left( p(\eta) + q(\eta) \right) - \frac{2\gamma}{\delta} \delta p'(\eta) \right] - (1-\phi)^{2.5} (1-\phi_2)^{2.5} M^2 q'(\eta) - Kq'(\eta) = 0
\]
\[
\left( \frac{\kappa_{nf}}{\kappa_f} + \frac{4}{3} \right) \theta''(\eta) + \left( 1-\phi_2 \right) \left( 1-\phi_2 \right) + \phi_2 \left( \frac{\rho C_p}{\rho_f} \right) \right] + \phi_2 \left( \frac{\rho C_p}{\rho_f} \right) \]
\[
\times Pr \left( p(\eta) + q(\eta) \right) \theta'(\eta) + Pr Ec \left( \left( p''(\eta) \right)^2 + \delta^2 \left( q''(\eta) \right)^2 \right) = 0
\]
\[
R^*(\eta) + \frac{Sc}{\left(1-\phi_2\right)} \left[ \left( p(\eta) + q(\eta) \right) R'(\eta) - R, R(\eta) \right] = 0
\]
The corresponding transformed boundary conditions are
\[
p = 0, \ \ p'(\eta) = 1 + \alpha p''(0), \ \ q = 0, \ \ q'(\eta) = \lambda + \alpha q''(0), \ \theta(\eta) = 1 + \beta \theta'(0), \ \ R(\eta) = 1 + \beta R'(0) \text{ at } \eta = 0
\]
\[
p' \to 0, \ \ q' \to 0, \ \theta \to 0, \ \ R \to 0 \text{ as } \eta \to \infty
\]
Non-dimensional quantities are
\[
\gamma = \frac{\Omega}{a}, \quad \alpha = k \sqrt{av_f}, \quad \beta = l \sqrt{a v_f}, \quad \lambda = \frac{b}{a}, \quad R = \frac{4 \sigma T^3}{\varphi \kappa_f}, \quad \delta = \frac{y}{x} = \sqrt{\frac{a y v_f}{x}}, \quad M = \sqrt{\frac{\sigma B_0^2}{a \rho_f}}
\]

(18)

\[
\text{Pr} = \frac{(\mu C_p) f}{\kappa_f}, \quad Ec = \frac{U_w^2}{(C_p) f (T_w - T_o)}, \quad K_1 = \frac{\mu_f}{a k \rho_f}, \quad S_C = \frac{v_f}{\beta_f}, \quad R_c = \frac{\xi (C - C_o)}{a}^{n-1}
\]

For mathematical and engineering point of concern, the skin friction coefficient, the reduced Nusselt number, and the Sherwood number are as follows

\[
C_{fs} = \frac{\mu_{lbf}}{\rho_f (ax)^2} \frac{\partial u}{\partial z} \bigg|_{z=0}, \quad C_{fs} = \frac{\mu_{lbf}}{\rho_f (ax)^2} \frac{\partial v}{\partial z} \bigg|_{z=0}, \quad N_{as} = \frac{x \kappa_{lbf}}{\kappa_f (T_w - T_o)} \frac{\partial T}{\partial z} \bigg|_{z=0}
\]

\[
S_{as} = -\frac{x}{(C_w - C_o)} \frac{\partial C}{\partial z} \bigg|_{z=0}
\]

\[
\text{Re}^2 C_{fs} = \frac{1}{(1 - \phi_1)^{2.5} (1 - \phi_2)^{2.5}} p''(0), \quad \text{Re}^{-1} \text{Re}^2 C_{fs} = \frac{1}{(1 - \phi_1)^{2.5} (1 - \phi_2)^{2.5}} q''(0),
\]

\[
\text{Re}^{-2} N_{as} = \left( \frac{\kappa_{lbf}}{\kappa_f} + \frac{4}{3} R \right) \theta'(0), \quad \text{Re}^{-2} S_{as} = -R'(0);
\]

Where \( \text{Re} = \frac{U_w x}{v_f} \) is the local Reynolds number.

3. Results and Discussion

The system of moderated coupled nonlinear ODE’s (13)-(16) including a new set of subject to condition (17) is utilized R-K-Fehlberg-integration with shooting method. The impacts of different governing dimensionless parameters are inspected, the Magnetic parameter \( M \), Porous medium parameter \( K_1 \), Rotation parameter \( \gamma \), Eckert number \( Ec \), Chemical reaction parameter \( R_c \), Schmidt number \( Sc \), Prandtl number \( Pr \), Thermal jump parameter \( \beta \), velocity slip parameter \( \alpha \), Thermal radiation parameter \( R \) on the velocities \( p'(\eta) \), \( q'(\eta) \), temperature \( \theta(\eta) \), and concentration \( R(\eta) \) fields, which are portrayed via Figs.2-17.

The performance of magnetic parameter \( M \) on stream components \( p'(\eta) \) and temperature dispersal \( \theta(\eta) \) is explained in Figs.2-3. With the influence of \( M \) in an electrically conducting liquid, drug-like energy is established. This resistive energy tends to slow the fluid movement at the expense of temperature improvement over the surface. This is often depicted by the decrease within the velocity values, which is found in Fig. 2. Because of a decrease in velocities components, an improvement happens in liquid temperature, as saw in Fig. 3.

The assessment between a base liquid (Ethylene glycol), a Cu- Ethylene glycol, and Cu- Fe3O4/ Ethylene glycol on transverse velocity and the degrees of hotness distributions is introduced in Figs. 4-5 Here, we see that compared to nanofluid (Cu-C2H6O2) and hybrid nanofluid (Cu- Fe3O4/ C2H6O2) of base fluid (C2H6O2). The velocity is high, but the opposite is true concerning temperature. It is found that the transverse velocity of the base liquid (Ethylene glycol) is more prominent as compared to the Cu-Ethylene glycol and Cu- Fe3O4/
C$_2$H$_6$O$_2$. However, the reverse trend is viewed in the temperature profile. The temperature for the Cu-$\text{Fe}_3\text{O}_4$/C$_2$H$_6$O$_2$ is greater than the C$_2$H$_6$O$_2$ and Cu-C$_2$H$_6$O$_2$. Physically, nanoparticles dissipate heat in the form of heat. Consideration of more nanoparticles can lead to more radiation that escalates the temperature.

Figs. 6-7 are drawn to investigate the effects of the velocity slip parameter $\alpha$ on the transverse velocity distributions of $p'(\eta)$, and $q'(\eta)$. It can be observed from these that the transverse velocity slip parameter $\alpha$ brings about a remarkable drop in the transverse velocity distribution $p'(\eta)$ and $q'(\eta)$ for both liquids. This is because, within sight of slip, the fluid's velocity near the surface doesn’t continue as correspondingly for the stretching surface velocity. Therefore, the slip velocity intensifies by escalates the velocity slip parameter $\alpha$. This is followed by reducing the liquid velocity because the deformation of the stretching surface, under the slip conditions, can only be transferred to the fluid.

The impact of the $\gamma$ over the transverse velocity and temperature profiles are presented in Fig. 8. From this figure, we can see that increasing the rotation parameters $\gamma$ significantly reduces the nanofluid's velocities (Cu-C$_2$H$_6$O$_2$) and the hybrid nanofluid (Cu-$\text{Fe}_3\text{O}_4$/C$_2$H$_6$O$_2$). Fig. 9. Shows the impact of permeable media on $p'(\eta)$. It is seen that the velocities of both nanofluids and hybrid nanofluids decrease as the parameter of the permeable media growth.

Fig. 10. It shows that the velocity escalates with the growth $\delta$. The Eckert number refers to the association between enthalpy and kinetic energy. It converts the conservation of kinetic energy into internal energy by work done against the viscous fluid stress. With the aid of action against viscous fluid tension, it transforms the conservation of kinetic energy into internal energy. The positive $Ec$ represents the diminishing concerning heat beside the sheet to the fluid. Fig. 11 shows that heat is transferred from the sheet to the nanofluid, in this manner increasing the temperature of the nanofluid. Fig. 12 shows the importance of the thermal slip parameter $\beta$ on the temperature field. The temperature profile is reduced by escalates the thermal slip parameter. Physically, less heat is transferred from the surface to the liquid as the thermal slip parameter grows, decreasing the temperature.

Fig. 13 represents the impact of thermal distribution $\theta(\eta)$ for different values of $R$. It is clear that the distribution of temperature at the boundary layer improves with varying values of the thermal radiation parameter. The effect of thermal radiation is to improve heat transfer because the thermal boundary layer increases with increasing thermal radiation. Therefore, it has been reported that the process of reducing thermal radiation should proceed at a faster rate. Fig. 14. This shows that the concentration field changes in Schmidt number hydrogen gasses, water vapor, and ammonia. It is found that the concentration field of carbon dioxide for hydrogen escalates regularly compared to water vapor. In this manner, water vapor can be utilized to maintain a concentration field, and hydrogen can be used to maintain a well-concentrated field. As the Schmidt number increasing, the concentration reduces with the thickness of the boundary layer. The concentration profiles for $Rc$ is shown in Fig. 15. It is seen that enhancing the values of chemical reaction reduced the liquid. As a result, the thickness of the concentration boundary layer escalates.
Figure 2. Influence of $P'(\eta)$ for $M$.

Figure 3. Influence of $\theta(\eta)$ for $M$.

Figure 4. Influence of $P'(\eta)$ for nanofluid.

Figure 5. Influence of $\theta(\eta)$ for nanofluid.

Figure 6. Influence of $P'(\eta)$ for $\alpha$.

Figure 7. Influence of $\theta(\eta)$ for $\alpha$.

Figure 8. Influence of $P'(\eta)$ for $\gamma$.

Figure 9. Influence of $P'(\eta)$ for $K_1$. 
Figure 10. Influence of $P'(\eta)$ for $\delta$.

Figure 11. Influence of $\theta(\eta)$ for $Ec$.

Figure 12. Influence of $\theta(\eta)$ for $\beta$.

Figure 13. Influence of $\theta(\eta)$ for $R$.

Figure 14. Influence of $R(\eta)$ for $Sc$.

Figure 15. Influence of $R(\eta)$ for $Re$.

Figure 16. $C_{f1} \cdot Re^2$ for $M$ for nanofluid and hybrid nanofluid in $y$ direction.

Figure 17. $N_{\mu} \cdot Re^{-1}$ for nanofluid and hybrid nanofluid in $y$ direction.
Fig. 16 elucidates the influence of $M$ and $\Omega$ on $\frac{1}{C_{fs}} \text{Re}^2$. It is noticed that skin friction enhances for higher values of $M$ while it increases via $\Omega$. Impact of the base fluid, nanofluid, and hybrid nanofluid on heat transfer presented in Fig. 17 the rate of heat transfer is less for base fluid as compared to the $\text{Cu} / C_2H_6O_2$ and $\text{Cu} – Fe_3O_4 / C_2H_6O_2$. Thus, the rate at which heat is higher transit is higher for hybrid nanofluids than the Ethylene glycol. Table 3 shows the numerical values of the friction factor for $M$, $\Omega$, $\gamma$, $\phi_s$.

| $M$ | $\Omega$ | $\gamma$ | $\phi_s$ | $\frac{1}{(1-\phi)^{\frac{1}{2}}}$ | $\frac{1}{(1-\phi)^{\frac{3}{2}}}$ | $\frac{1}{(1-\phi)^{\frac{3}{2}}}$ | $\frac{1}{(1-\phi)^{\frac{3}{2}}}$ | $\frac{1}{(1-\phi)^{\frac{3}{2}}}$ |
|-----|----------|----------|----------|----------------------------------|----------------------------------|----------------------------------|----------------------------------|----------------------------------|
| 0.0 | 0.1      | 0.5      | 0.05     | 2.14464                          | 2.28453                          | 1.28679                          | 1.48516                          |
| 0.5 | 1.0      |          |          | 1.84580                          | 2.14121                          | 1.54525                          | 1.58270                          |
| 1.0 |          |          |          | 1.94614                          | 2.36551                          | 1.63747                          | 1.68994                          |
| 1.5 | 0.2      |          |          | 2.41543                          | 2.83606                          | 1.96505                          | 2.27353                          |
|     | 0.3      |          |          | 1.85341                          | 2.30152                          | 1.43147                          | 1.89894                          |
|     | 0.5      |          |          | 1.88474                          | 2.31105                          | 1.96515                          | 2.28078                          |
| 0.2 |          | 0.2      |          | 1.99092                          | 2.42500                          | 0.92152                          | 1.16418                          |
|     | 0.4      |          |          | 1.89593                          | 2.43401                          | 0.81552                          | 1.15529                          |
| 0.6 |          | 0.6      |          | 1.94542                          | 2.32153                          | 1.43346                          | 1.67775                          |
|     | 0.04     |          |          | 1.77374                          | 2.2006                          | 2.67838                          | 2.01750                          |
| 0.08| 2.15620  |          |          | 2.21470                          | 1.51564                          | 1.76095                          |
| 0.1 |          |          |          | 2.28421                          | 2.92372                          | 1.97656                          | 2.50133                          |

We observed that, the skin friction coefficient enhances for intensified values of the magnetic parameter, Angular velocity, rotation parameter $\phi_s$.

4. Conclusions

The convective heat transfer and hybrid nanofluid flow over a sheet with slip impacts. Hybrid nanofluid is considered by suspending two different nanoparticles, which are Cu and then Fe$_3$O$_4$ is C$_2$H$_6$O$_2$. When utilizing the similarity transformations, the overseeing conditions with boundary conditions must be transformed into a system of ODEs, discussed graphically. The main conclusions are summarized as follows: It is noticed that the hybrid nanofluid Cu–Fe$_3$O$_4$/C$_2$H$_6$O$_2$ is best as a heater on growing $M$. Conversely, on increasing $Ec$, the slip parameters and stretching, it is higher as a cooler (warmer) in contrasting Fe$_3$O$_4$/C$_2$H$_6$O$_2$ (Cu/C$_2$H$_6$O$_2$); The hotness of the nanofluid grows with increasing of Eckert number; A decrease in the fluid velocity is observed for higher values of the velocity slip parameter; The transverse velocity profile reduces with increasing solid volume fraction and porous medium but a reverse trend is observed with the increments of $\gamma$, $K_1$, $\alpha$ parameters; The nanofluid velocity decreases as the existence of the $M$ becomes stronger; The rate at which heat transfer occurs for Cu/C$_2$H$_6$O$_2$ is the highest among Cu–Fe$_3$O$_4$/C$_2$H$_6$O$_2$ nanofluids.

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Conflicts of Interest

The authors declare no conflict of interest.

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