Contactless temperature measurement in wire-based electron beam additive manufacturing Ti-6Al-4V

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Abstract

The complex thermal cycles and temperature distributions observed in additive manufacturing (AM) are of particular interest as these define the microstructure and the associated properties of the part being built. Due to the intrinsic, layer-by-layer material stacking performed, contact methods to measure temperature are not suitable, and contactless methods need to be considered. Contactless infrared irradiation techniques were applied by carrying out thermal imaging and point measurement methods using pyrometers to determine the spatial and temporal temperature distribution in wire-based electron beam AM. Due to the vacuum, additional challenges such as element evaporation must be overcome and additional shielding measures were taken to avoid interference with the contactless techniques. The emissivities were calibrated by thermocouple readings and geometric boundary conditions. Thermal cycles and temperature profiles were recorded during deposition; the temperature gradients are described and the associated temperature transients are derived. In the temperature range of the $\alpha+\beta$ field, the cooling rates fall within the range of 180 to 350 °C/s, and the microstructural characterisation indicates an associated expected transformation of $\beta \rightarrow \alpha'+\alpha$ with corresponding cooling rates. Fine acicular $\alpha$ and $\alpha'$ formed and local misorientation was observed within $\alpha$ as a result of the temperature gradient and the formation of the $\alpha'$.

Keywords Ti-6Al-4V · Additive manufacturing · Contactless temperature measurement · Wire-based

1 Introduction

Additive manufacturing (AM) is an emerging technology that offers a variety of new and interesting possibilities. Using this technology, companies can produce finished and semi-finished products that have optimised geometry, technological and metallurgical properties but also economic benefits. AM covers a wide range of processes, which can generally be classified according to the ISO 17296-2:2015 standard [1] as follows: vat photopolymerisation, material jetting, binder jetting, powder bed fusion (PBF), material extrusion, directed energy deposition (DED) and sheet lamination. The PBF method is widely used to process metallic materials, but economical DED processes (e.g. LDM, WAAM) are also suitable for the production of AM components. Numerous metallic materials can be used, such as steels [2–4], aluminium materials [5–7] and nickel-based alloys [8–10]. Titanium and its alloys are of particular interest, especially to the aerospace industry, due to their excellent properties, including their high strength to density ratio and their resistance to corrosive environments [11]. However, the route followed to produce such alloys is highly complex, cost-intensive and has considerable drawbacks. To overcome these disadvantages, a high material utilisation ratio should be achieved (buy-to-fly ratio); this can be achieved by using DED AM processes [12]. When reactive materials are used in production, such as Ti-based alloys, few suitable DED processes are available, and a remarkable number of requirements must be met to create an appropriate shielding environment to avoid contamination. To prevent the oxidation of the AM parts produced using the WAAM process, it is necessary to provide a high argon gas flow rate. This, in turn, leads to significant forced convection cooling, and careful attention needs to be paid to thermal management and interpass conditions [13–15]. Oxidation and dissolved
oxygen in the titanium matrix causes an undesirably rapid reduction in the fracture elongation and increases its strength due to solid-solution hardening [16, 17]. Electron beam processes serve as a useful alternative; because the process is carried out under vacuum, they are suitable for processing titanium and contamination by oxygen is avoided [18]. Electron beam freeform fabrication (EBF$^3$) is one such wire-based DED process that uses an electron beam as a heat source. The feasibility of applying this process to titanium alloys has been already reported in the literature [19–24].

All AM processes display complex multiple thermal cycles and transient temperature profiles. The temperature distribution influences many factors, such as the melt pool geometry [25–27]. The thermal history of the components defines various metallurgical phenomena, microstructural features, and properties related to these components [11, 28]. In addition, by taking temperature measurements and gaining knowledge about temperature profiles, on-line process control / monitoring can be facilitated and applied for thermal non-destructive evaluation (NDE) purposes to stabilize and further improve the process [29–31]. Consequently, researchers have shown particular interest in increasing the knowledge about the thermal cycles and temperature distribution in AM by conducting experimental measurements and numerical thermal simulations. Due to the nature of the process, which involves intrinsic layer-by-layer material stacking and continuous build-up, contact methods to measure temperature are not as suitable as contactless methods. The use of the latter, such as thermal imaging or point measurement methods, is required to collect spatial and temporal temperature information for the region of interest [32].

Li et al. [31] summarised the recent trends and provided an overview of possible contactless in-situ temperature measurement techniques that could be applied in metallic AM. These contactless techniques include conventional infrared irradiation temperature measurement methods (single-wavelength / one-colour technique) and two- or more wavelength (multicolour-) techniques [31]. When multi-colour pyrometers are applied, it is generally assumed that the ratio between the emissivities for the respective spectral ranges is known and that the emissivity ratios are constant, rather than temperature-dependent. Each material has its own emissivity-to-wavelength relationship, and costly multi-colour pyrometers with a specific wavelength range can often not be used to measure temperature with sufficient accuracy due to the changing emissivity ratio of the material [31–33]. One-colour infrared irradiation techniques offer cost-effective alternatives; these require the user to have knowledge about and quantitative values for the emissivity of a material as well as its present condition [32]. The material emissivity needs to be carefully calibrated and selected before such measurements are taken, as the emissivity varies according to the wavelength, temperature and surface conditions. The latter include the roughness, contamination level, oxidation level or surface stress state of the material [34, 35]. Due to its simplicity, robustness and cost efficiency, this contactless technology has already been tested in WAAM processes to investigate the effect of heat accumulation on the interlayer temperatures [33, 36, 37], alterations in arc characteristics [38] and changes in metallurgical and mechanical properties [36, 37, 39] during the layer-by-layer stacking process. However, several points need to be considered when WAAM processes are carried out, such as the effects of transient surface contamination (e.g. oxidation, slag and smoke), changing transmissivity/optical path (welding fume), and interference from the arc light, since the arc produces radiation that covers the entire spectrum. In an electron beam freeform fabrication process, on the other hand, the presence of vacuum prevents surface contamination (which can otherwise be comparatively severe), provides a more constant transmittance/optical path and prevents interference from a heat source. On the other hand, new challenges also arise when performing contactless temperature measurements during metal processing under vacuum. Limitations in accessibility and devices that are suitable for use under vacuum conditions need to be taken into account. Vacuum feedthroughs for electrical signals are limited and the resulting X-rays during beam exposure may damage the electronics implemented in the chamber without additional shielding measures. During processing and metal deposition, elements with a high saturated vapour pressure tend to evaporate under vacuum conditions and can interfere with contactless measurements. The evaporative loss of aluminium that occurs when Ti-6Al-4V is processed under vacuum has already been reported in the literature [21, 23, 40–42]. Additional measures have to be taken to ensure that the temperature can be recorded continuously and without disturbance inside the vacuum chamber. Contactless measurements on workpieces inside the chamber, e.g. by IR pyrometers, can in principle also be carried out from outside the chamber [43]. Electron beam vacuum chambers are designed with a viewing port to provide a direct view into the chamber through a small (lead) glass window. However, measurements from outside the chamber are subject to uncertainties because the lead glass used in the viewing port provides sufficient optical protection from UV and IR radiation, changes in transmissivity, and an increase in the measurement distance between the pyrometer and the region of interest, resulting in low spatial resolution.

The recent literature [30] contains little information about contactless temperature measurements in vacuum-based EBF$^3$. The present study was carried out to demonstrate the feasibility of contactless temperature measurement methods when they are implemented inside a vacuum chamber (pyrometer measurement method) and under normal atmospheric conditions (thermal imaging technique). Appropriate methods for positioning and protecting the devices from the metal vapour during processing are presented. Thermal cycles and
temperature profiles, described by presenting isotherms, are determined during electron beam AM, and the cooling rates are associated with the observed microstructure of the deposited weld metal.

2 Experimental

2.1 Materials

A solid wire of titanium alloy Ti-6Al-4V AWS A5.16 ER Ti5 (EN ISO 24034) with a diameter of Ø 1.2 mm was used in this study. All experiments and the weld bead deposition were performed on a titanium grade 5 substrate with dimensions of $200 \times 50 \times 13$ mm. The chemical composition of the applied materials is listed in Table 1. The proportion of the elements aluminium (Al), vanadium (V) and iron (Fe) were determined by atomic absorption spectroscopy. The interstitial element carbon (C) was identified by applying the solid-state infrared absorption detection method with LECO CS230 (LECO Corporation, St. Joseph, MO, USA) and nitrogen (N) was identified by applying the thermal conductivity detection method with LECO ON 736 (LECO Corporation, St. Joseph, MO, USA).

2.2 Equipment

2.2.1 Electron beam welding device

The experiments were performed using the electron beam welding device pro-beam EBG 45-150 k14 (pro-beam GmbH & Co. KGaA, Gilching, Germany). A maximum acceleration voltage and beam current of up to 150 kV and 300 mA can be supplied by this device. The working chamber provides a volume of approx. $1.4 \text{ m}^3$ with a vacuum level below $5 \times 10^{-3}$ mbar. Inside the chamber, workpiece handling was performed by manipulating a three-axis work table, while the electron beam gun and welding position PA/1F were fixed. The solid wire was fed by pressure rollers and ended in a water-cooled wire nozzle which could be moved along all three axes. The angle between the substrate and the fed wire was $35^\circ$ [44], and the distance between the wire nozzle and the substrate surface was set at 8 mm. A height distance of 1 mm between the wire tip and the substrate surface was set to enable a stable liquid metal bridge material transfer to take place during the process.

2.2.2 One-colour pyrometer

To perform pyrometry, a one-colour pyrometer Sensortherm METIS M318 (Sensortherm GmbH, Sulzbach, Germany) was used with a temperature range and spectral range of 150 to 1200 °C and 1.65–2.1 μm. The emissivity and transmissivity can be selected independently and values can be set variably with this device. The pyrometer was mounted on the wire feed unit in the vacuum chamber and aligned at $35^\circ$ at a radial distance of approx. 450 mm from the substrate. The focal point was adjusted on the material surface, and a spot size of Ø 2 mm was used in all experiments.

When Ti-6Al-4V is processed under a vacuum, mainly aluminium evaporates from the melt pool. These evaporation losses coat the pyrometer lens, resulting in continuous changes in the transmittance and interference with the contactless measurement. To keep this evaporation from influencing the measurement, the setup was extended and a protective device was designed and attached to the pyrometer. The device operated as follows: a heat-resistant polyimide film (up to 400 °C) is moved continuously in front of the pyrometer lens (Fig. 1a) by a vacuum-compatible electric motor. In this way, the thickness of the aluminium coating is kept constant and depends mainly on the speed with which the film is moved. If the conveying speed is fast enough, the aluminium condensation on the polyimide film is no longer visible macroscopically, and it is possible to continuously record the temperature (Fig. 1b).

2.2.3 Near-infrared camera

Temperature fields were recorded with a near-infrared (NIR) camera Optris PI 1M (Optris GmbH, Berlin, Germany) that covered the spectral range of 0.85–1.1 μm and a temperature range of 675–1800 °C. The NIR camera produces images with a resolution of $764 \times 480$ pixels and has a maximum frame rate of 32 Hz. The emissivity and transmissivity can be selected independently and values can be set variably with this device. The camera was positioned outside the vacuum chamber in the optical viewing system in addition to the integrated CCD camera. A dichroic mirror acts as a wavelength-specific beam splitter and enables the use of the CCD camera to

| Table 1 | Measured chemical composition of applied filler wire AWS A5.16 ER Ti5 and Ti-6Al-4V substrate; n.m., not measured |
|---------|---------------------------------------------------------------|
|         | Al (wt. %) | V (wt. %) | Fe (wt. %) | Ti (wt. %) | C (wt. %) | N (wt. %) | O (wt. %) |
| Filler wire | 6.36    | 3.48    | 0.11    | bal.       | 0.018    | < 0.005  | n.m.      |
| Substrate  | 6.38    | 3.49    | 0.21    | bal.       | 0.015    | < 0.005  | n.m.      |
monitor the melt pool at the same time that the NIR camera is used to record the corresponding temperature field. The permanently installed foil transport system in front of the optical viewing prevents a coating of the CCD camera and also of the NIR camera. The schematic illustration of the experimental setup is depicted in Fig. 2.

2.3 Experimental setup

2.3.1 Determination of the emissivity

The one-colour pyrometer covers a wide temperature range from above the $\beta$-transus temperature to low interlayer temperatures, and a suitable emissivity needs to be determined. The emissivity $\varepsilon_P$ was calibrated for the pyrometer system using a benchmark test in which the thermal cycles recorded with the pyrometer and the thermocouples were simultaneously opposed in the temperature range of approx. 200–1000 °C. In the benchmark test, the Ti-6Al-4V sheet was continuously heated by exposing it to an electron beam for 240 s with an adapted beam figure of concentric circles (maximum diameter of 8 mm; the other diameters as follows: 1.8, 3.8, 5.4, 6.6, 7.4 and 7.8 mm) (Fig. 3a). Electron beam processing is characterised by the steep temperature gradient observed from the beam spot to the surrounding area. As such a steep temperature gradient is not desirable during the calibration, the electron beam was slightly defocused to reduce the temperature gradient. The applied process parameters are listed in Table 2.

Three thermocouples (K1, K2, K3) of type K with diameters of Ø 0.2 mm were spot-welded along the diameter of a circle that was Ø9 mm in the circumferential direction. The reduced diameter of the thermocouples combined with spot welding the thermocouples to the workpiece results in a reduction in response lag [45]. The recorded temperature cycles were then compared with those recorded with a contactless single-colour pyrometer, applying different emission values when the measurement spot was placed at the same radial distance from the centre. A suitable emission coefficient was selected on the basis of a minimum relative temperature difference for the temperature range of interest in the heating and cooling stages.

2.3.2 Contactless temperature measurements in electron beam additive manufacturing

Single-track weld beads of Ti-6Al-4V, each with a length of 50 mm, were deposited on the substrate. The relevant main process parameters and beam oscillation parameters are listed in Table 3. As beam figure, ten concentric circles with...
diameters from 4.0 to 0.4 mm in steps of 0.4 mm were used. The suitability of the parameters has already been demonstrated in [23], and the selected parameter configuration ensures a proper material transfer by the preferred liquid metal bridge transfer mode. Continuous material transfer from the filler wire to the molten pool via the liquid metal bridge results in a smooth and uniform weld surface. The material deposition with the selected parameter configuration results in a track width of about 5.6 mm and an average height of about 1.7 mm [23].

The temperature during EBAM at a nominal constant distance of 10 to 35 mm (5 mm steps from the beam centre on the weld seam centreline) was recorded with a pyrometer using a previously calibrated emissivity $\varepsilon_P$ (Fig. 4). The recorded temperature curves were evaluated to identify the peak temperatures in the quasi-stationary stage as a function of their relative offset from the beam centre. This evaluation allowed the derivation of the temperature field, the associated gradients and the cooling rates. The experimental setup is schematically depicted in Fig. 4.

| Unit       | Values |
|------------|--------|
| Acceleration voltage | $U_{acc}$ kV 90 |
| Beam current | $I_{beam}$ mA 8 |
| Heating time | $t$ sec 240 |
| Focal point | $f_p$ - Substrate surface and defocused by +100 mA lens current |
| Working distance | - mm 800 |
| Beam figure (bop) | - Concentric circles |
| Frequency (bop) | $f$ Hz 1000 |
| Amplitude of deflection (bop) | $x,y$ mm Ø 8 |

In addition, the temperature field was recorded using the thermal imaging technique and a near-infrared camera. The images were acquired at a frame rate of 8 Hz for the single-track deposition of Ti-6Al-4V, and the region of interest (ROI) was exported at a resolution of 140 × 60 pixels. The data were analysed by using MATLAB R2017a software (MathWorks, MA, USA).

### 2.4 Metallographic characterisation

In order to correlate the derived cooling rates with the microstructure of the material, the corresponding cross section was examined. The deposited Ti-6Al-4V was cut perpendicular to the welding direction, and the resulting cross section was hot mounted, ground using SiC paper, and polished using a silica oxide polishing suspension (OPS) solution prior to the microstructural analysis. The scanning electron microscopy (SEM) investigation was performed using a TESCAN Mira3 microscope (TESCAN, Brno, Czech Republic) equipped with a Hikari detector and a TSL-OIM Data Collector software package for electron backscatter diffraction (EBSD) analysis. EBSD measurements were performed on an area of 50 $\mu$m × 50 $\mu$m, with a step size of 0.1 $\mu$m. An acceleration voltage of 15 kV, working distance of 15 mm and spot size of 8 nm were used during the EBSD measurements. The generally accepted transition angle between LAGB and HAGB is 15° in the case of an ideal arrangement of dislocations [46]; this was used to define a high-angle grain boundary. Using the OIM Analysis v.8 software, a minimum grain size of 0.5 $\mu$m was chosen to guarantee that each grain consisted of at least five measured pixels. The confidence index was standardised. The EBSD data was then cleaned, considering a minimum confidence index of 0.1 correlated to the neighbour grains.

![Beam Figure](image)

**Fig. 3** Schematic illustration: a concentric circles beam figure and b experimental setup used to calibrate emission coefficient by comparing thermal cycles recorded by thermocouples and one-colour pyrometer
3 Results and discussion

3.1 Determination of the emissivity

According to the literature [35, 47, 48] and preliminary investigations, an emissivity in the range of 0.280–0.325 proved to be most promising for calibrating the emissivity for the system and the material surface condition. Depending on the set emissivity, the temperature curves recorded using the pyrometer without contact (solid line) differed from those recorded with contact with thermocouples (dashed line) (Fig. 5).

It was observed that the consistency of the temperature cycles depended not only on the emissivity, but also whether this emissivity was measured during the heating or cooling stage. In the initial phase of heating (0–5 s), a maximum deviation in the contact and contactless curve could be observed for all tests carried out, regardless of the emissivity chosen. The respective deviations can be attributed to the different response times, spatial resolution offered by the applied methods, which result from the different spot sizes (pyrometer Ø 2 mm, spot-welded thermocouples Ø 0.2 mm), as well as differences in the existing transient temperature field and the resulting steep temperature gradients. Since the measured spot sizes and, thus, the range over which the temperature is averaged differed, the average temperature levels as measured by the pyrometer and thermocouples also differed significantly (Fig. 6).

Table 3 Summary of input and beam oscillation parameters (bop) for single-track Ti-6Al-4V deposition

| Unit          | Values |
|---------------|--------|
| Acceleration voltage \( U_{\text{acc}} \) | kV 90 |
| Beam current \( I_{\text{beam}} \) | mA 17.5 |
| Welding speed \( v_{\text{weld}} \) | mm/s 9.0 |
| Wire feed rate \( V_{\text{wire}} \) | m/min 3.3 |
| Feed angle \( \alpha_{\text{feed}} \) | ° 55 |
| Focal point \( f_p \) | - Substrate surface |
| Beam figure (bop) | - Concentric circles |
| Frequency (bop) \( f \) | Hz 1000 |
| Amplitude of deflection (bop) \( x_{\text{def}} \) | mm Ø 4 |

As the heating continued, and the temperature rose, the transient temperature gradient flattened, leading to an improved approximation of the curves measured. The offset between the temperature levels could be adjusted by changing the emissivity, and the selected emissivity \( \varepsilon_p \) of 0.300 and 0.325 resulted in an acceptable alignment of the curves for the temperature range of interest. Ending the heating stage after 240 s by switching off the beam initiated the cooling stage. Since the sample is exposed to vacuum, the heat slowly dissipates due to thermal radiation and heat conduction without the contribution of convection.

The deviation between the contactless and contact measurement values observed during the benchmark test can be described by the relative and absolute difference related to the temperature (Fig. 7). In the initial stage of the process at a lower temperature (200–400 °C), i.e. a rapidly evolving transient temperature field, contactless measurement is characterised by a high degree of relative difference and a temperature deviation of 45–15% and 180–70 °C, respectively. As the exposure time to elevated temperature increases, the temperature transient softens, indicative of an almost steady state. This implies that an equilibrium has occurred between the energy input introduced by the beam and the heat that dissipates by radiation and heat conduction. In each respective stage, the relative and absolute difference can be minimised by setting an appropriate emissivity level. At a given emissivity \( \varepsilon_p \) of 0.325, the relative and absolute difference at the end of the heating stage (minimum temperature gradient) and the subsequent cooling stage (moderate temperature gradient) can be reduced to almost zero. Therefore, an emissivity \( \varepsilon_p \) of 0.325 was selected for the one-level pyrometer for further experiments.

3.2 Contactless temperature measurements in electron beam additive manufacturing

3.2.1 Thermal cycles by pyrometry

Thermal cycles recorded by pyrometry with an emissivity \( \varepsilon_p \) of 0.325 at a distance of 10 to 35 mm from the beam centre are shown in Fig. 8. Before the actual metal deposition occurs and the process starts (< 0 s), the initial temperature of 150 °C
represents the lower threshold temperature of the pyrometer. This means that only temperatures that exceed this limit can be measured. After the process is initiated and the electron beam is switched on, the temperature increases. The first maximum can be considered as the substrate temperature, since the position of the pyrometer measuring spot is still located on the substrate. Due to the movement of the machine table/substrate during the process, the measuring spot shifts from the substrate to the recently deposited weld bead and records the surface temperature of the deposited metal. Therefore, the second temperature increase represents the transition of the measuring spot from the heated substrate to the deposited weld metal. The maximum temperature is reached at the quasi steady-state temperature. The peak temperatures are dependent on the actual distance of the measuring spot from the beam centre, and these are nearly constant over the entire process duration (Fig. 8a). At the end of the process, the metal deposition stops, and the electron beam is switched off, resulting in a drop in temperature due to cooling.

As the initial position of the measuring spot of the pyrometer in relation to the centre of the beam can only be set manually, the nominal relative offset is subject to uncertainties. The relative positions need to be corrected, which is performed by means of applying the time derivative for the temperature curves (Fig. 8b). The actual relative position is determined by taking into account the period of time \( \Delta t \) between the first peak at process initiation \( t \)}
and the second peak at the transition from the substrate to the weld metal. Considering the predefined welding speed $v_w$ (9 mm/s), the actual relative position can be derived and calculated as follows:

$$x_{pos} = v_w \times \Delta t$$

Based on the pyrometer measurements and the recorded quasi steady-state temperatures for corrected relative beam offset positions, the temperature profile was interpolated (Fig. 9a) and the associated cooling rates were derived (Fig. 9b). In the temperature range of 1000–850 °C ($\alpha+\beta$ field), maximum and a minimum cooling rates of approx. 350 and 220 °C/s were determined. For the lower temperature level of 600 °C, a cooling rate of about 80 °C/s was derived.

3.2.2 Thermal cycles by thermal imaging

In addition to the pyrometer measurements, the temperature profile was recorded with a near-infrared camera during the single-track deposition. Since the thermal cycles were exported for different emissivities in the range of 0.1 to 1.0, numerous temperature profiles in the $y$- (Fig. 10b), but also in the $x$-direction (Fig. 10c) are possible. As with the one-colour pyrometer, a suitable emissivity needs to be selected. To determine a suitable emission coefficient, the temperature profiles for the $x$- and $y$- directions were evaluated with regard to the following geometric boundary conditions:

1. For the applied parameter configuration, the welding bead width is known and is approx. 5.6 mm [23]. Since the weld bead is symmetrical and the width is known, it can be assumed that the temperature recorded in the width of the weld bead should be equal to the solidus temperature at a certain emissivity ($T_{\text{solidus}} \approx T(0,w/2,\varepsilon)$) (Fig. 10b).

2. Tang et al. [27] simulated the transient fluid flow of the weld pool in EBF3, showing that the high temperature range is restricted to the area of direct electron beam exposure and has a steep temperature gradient to its surroundings. In the present study, the beam was continuously deflected to a
A circular area with a diameter of 4 mm. Therefore, in the outer area of direct beam exposure, a temperature above the maximum recordable temperature of the thermal imaging camera \(T_{\text{limit}} = 1800 \, ^\circ\text{C}\) must be expected \(T_{\text{limit}} < T(0,\varnothing_{\text{Beam}},\varepsilon)(\text{Fig. 10b})\). Considering the limitations, only the temperature profile of the area with non-direct beam exposure can be captured.

Taking into account the temperature profile in the \(x\)-direction, a change in the temperature gradient during cooling can be noted (Fig. 10c). The near-infrared camera covers the temperature range from 1800 to 675 \(^\circ\text{C}\), and the surface temperature is not only measured in the solid but also in the semi-solid and liquid states. The thermophysical properties (i.e., heat capacity, thermal conductivity, emissivity) vary with the temperature and physical state; heat flux/temperature profile changes noticeably at these transitions. The changes in the temperature gradients that occur during the cooling process can be assigned to the transitions from the liquid to the semi-solid and to the solid state, and the corresponding liquidus and solidus temperatures are known for Ti-6Al-4V (Fig. 10c).

Based on the temperature profiles in the \(x\)- and \(y\)-directions and the presented boundary conditions, the near-infrared camera was calibrated with an emissivity \(\varepsilon_T\) of 0.3. This value is close to the value used for the pyrometer, as the spectral ranges approach each other, and it correlates with values cited in the literature for titanium and titanium alloys [35, 47, 48] with similar spectral ranges measured under a vacuum.

Figure 11a shows the temperature profile as described by its isotherms in the quasi steady-state stage with a calibrated emissivity of 0.3. The melt pool (liquid) and mushy zone (liquid + solid) length is approx. 6.5 mm and 2.1 mm, respectively. Temperatures above the \(\beta\)-transus can be observed at a distance of up to 15 mm from the beam centre. The associated cooling profile was also derived (Fig. 11b), estimating rapid cooling rates (up to 1000 \(^\circ\text{C/s}\)) in the area just beyond the heat source.

Based on thermographic measurements, the temperature profile was interpolated (Fig. 12a) and the corresponding cooling rates were derived (Fig. 12b). In the temperature range of 850–1000 \(^\circ\text{C}\), slower cooling rates of approx. 180–350 \(^\circ\text{C/s}\) were obtained.

3.2.3 Comparison recorded thermal cycles by thermal imaging to pyrometry

Figure 13 depicts the interpolated temperature curves as determined by using both contactless methods. While the thermal imaging camera covers the higher temperature levels, the pyrometer mainly covers temperature levels lower than 1100 \(^\circ\text{C}\). A superimposed temperature range for both methods is limited to the range of about 800–1100 \(^\circ\text{C}\). For temperatures below 875 \(^\circ\text{C}\), a significant deviation in the curves was observed; measurements made with the thermal imaging camera were higher than the surface temperature recorded by the pyrometer. This may be attributed to the fact that this temperature level recorded by pyrometer is already within the range of the lower temperature threshold that can be measured with the thermal imaging camera. The low intensity thermal radiation (reflections) reflected from the ambient zone inside the chamber may indicate a higher temperature level for the region of interest than is actually present. The presence of thermal reflections, especially when conducting experiments with metals, are a well-known drawback in thermal imaging, which needs to be taken into account when analysing the results. Other aspects also need to be considered, such as the selection of the respective constant emissivity for the respective devices. While the emissivity for the one-level pyrometer was selected based on a calibration for the temperature range 400–1000 \(^\circ\text{C}\), the emissivity for the NIR camera was calibrated based on temperatures from 1600 \(^\circ\text{C}\) upwards. Since the overlapping temperature range is not in the primary calibration range of both

![Fig. 9](https://example.com/fig9.jpg)

(a) Contactless measured quasi steady-state temperatures depending on the actual relative position and derived exponential fit for temperature distribution.

(b) Derived associated function of the cooling rate
techniques at once, a degree of uncertainty in the accuracy of this range can be assumed.

3.3 Microstructural characterisation

The microstructure of Ti-6Al-4V is sensitive to the thermal history, and cooling conditions play an essential role in the formation of the different morphological features of the \( \alpha \)-phase [49]. Transition products from the \( \beta \) decomposition can arise and depend mainly on the thermal history of the deposited material. The evolving microstructure can consist of primary, secondary, plate-like, colony, grain boundary, acicular \( \alpha \), Widmanstätten structures, or hexagonal \( \alpha' \) (martensite) [11, 50].

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Fig. 10  a Recorded temperature field with an emissivity coefficient of \( \varepsilon = 0.3 \) and temperature profiles (taken along solid white lines) in b \( y \)-direction \( T = T(y, \varepsilon) \) and c \( x \)-direction \( T = T(x, \varepsilon) \) for variable emissivity ranging from \( \varepsilon = 0.1–1.0 \) for a single-track deposition in the 1st layer.
Ahmed et al. [51] described the phase transformations that occurred during the cooling process in α + β alloy Ti-6Al-4V and summarised the results in a schematic diagram that depicted a continuous cooling event. While cooling rates of higher than 410 °C/s lead to the formation of a fully martensitic α' microstructure, lower cooling rates result in the formation of a mixed microstructure consisting of α' and different morphological forms of diffusion-controlled α (Fig. 14) [51].

The microstructure of the single-track deposition transverse to the welding direction was investigated using a BSE detector and EBSD, which is shown in Figs. 15 and 16, respectively. The different orientations of the lamellas of α/α’ are shown in the IPF map in Fig. 16a. Their maximum thickness is ~ 1 μm. Zones of very fine and undistinguishable α/α’ are shown between the thicker α phases, as highlighted by the yellow dashed circles in Fig. 15.
The different sizes of the platelets of α or acicular shape of α’ are shown in Fig. 16b. The local misorientation distribution is illustrated by the kernel average misorientation (KAM) map in Fig. 16c. Finally, Fig. 16d shows the distribution of the confidence index for the investigated area. The very fine particles shown in Fig. 16c also exhibit a confidence index (CI) in Fig. 16d, indicating that they are α’. The thicker particles with higher image quality (IQ) values and low KAM are considered to be fine platelets of fine α-phase. If the regions in blue and green in Fig. 16d are considered as representative of α’, the fraction of α’ can be estimated to be ~40%, while the remaining 60% would correspond to acicular α-phase. The complex thermocycle with a rapid cooling period leads to complex residual stress distribution. This, in turn, causes local plastic deformation and leads to the formation of a local misorientation spread and substructures within the α/α’. The massive formation of α’ and the temperature gradient within the single track that occurs during the rapid cooling phase serve as the sources of the residual stresses. The regions indicated by the black arrow in Fig. 16c are regions where local plastic deformation occurred during the formation of the microstructure in the deposition phase. The measurements refer only to the first single track. Subsequent layers further reduce the stresses, as observed in [23].

4 Summary and conclusion

In the present work, it was possible to demonstrate the feasibility of applying contactless infrared irradiation temperature measurement in electron beam additive manufacturing. Two different techniques, pyrometry and thermal imaging, were used to collect spatial and temporal information while metal was deposited on a substrate under vacuum. The near-infrared camera could be used to cover the higher temperature range (i.e. the transition from liquidus to semi-solid and on to solid), while the one-colour pyrometer could mainly be used to record lower temperatures (i.e. at which diffusion and structural changes take place). Based on the recorded and derived temperature profiles, predictions for the evolved microstructure were made and were compared with the characterised weld metal deposit. The results can be summarised as follows:

4.1 Pyrometry

- The evaporation loss of aluminium that occurred when Ti-6Al-4V is processed interferes with the contactless temperature measurements made in the vacuum chamber. Placing a continuous foil transport in front of the pyrom-
A benchmark test, including heating and cooling stages, was designed to calibrate the pyrometer with thermocouples and compare the recorded thermal cycles. In the initial stage of heating, the associated steep temperature gradient results in comparatively high degrees of relative difference and deviation due to the different spatial resolution of the spot-welded thermocouples and pyrometer. At later stages of heating and in the cooling stage, the lower and more moderate temperature gradient results in more accurate results and higher consistency among the compared temperature curves.

Regarding the spectral range covered by the single-colour pyrometer and the surface condition of the material tested, a calibrated emissivity of 0.325 results in a minimal relative difference for the temperature region of interest and in the cooling stage. The corresponding emissivity for the one-colour pyrometer was chosen for the subsequent experiment and metal deposition.

Single-track metal deposition of Ti-6Al-4V was performed, and temperature cycles were recorded for several positions relative to the beam centre. To determine the spatial temperature distribution, the quasi-stationary peak temperatures were evaluated as a function of the offset from the beam centre and mathematically described by exponential fitting in order to examine the temperature gradient and derive the corresponding temperature transient function.

- Based on the measurements and interpolations, cooling rates of approx. 220–350 °C/s could be determined by contactless pyrometry in the temperature range of 850–1000 °C (α+β field). At lower temperatures, e.g. 600 °C, a cooling rate of about 80 °C/s is still possible.

### 4.2 Thermal imaging

- Temperature profiles for emissivities between 0.1 and 1.0 for the quasi steady-stage stage during weld metal deposition were exported. Geometric boundary conditions and change in temperature gradient were used to calibrate an emissivity of 0.3.
- The temperature profile for an emissivity of 0.3 was exported, and the associated temperature transient profile was derived. In addition, a phase transformation starts to take place in the rear area and 15 mm away from the beam centre. Locally, high cooling rates (up to 1000 °C/s) were measured in the vicinity of the direct beam exposure, and a decrease was observed towards the rear as the distance from the electron beam and metal deposition increased.
- Based on the measurements and interpolation, it was possible to record cooling rates of approx. 180–350 °C/s by...
contactless thermal imaging in the temperature range of 850–1000 °C (α+β field).

4.3 Microstructure characterisation

- A microstructure consisted of fine acicular α phase and α' was observed after a single-track deposition.
- The residual stress distribution in the very first single track, which caused local plastic deformation and resulted in the formation of local misorientation during the cooling phase, can be characterised by EBSD measurement.

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Declarations

Conflict of interest The authors declare that they have no conflict of interest.

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