Study on Micro-Crack Induced Precision Severing of Quartz Glass Chips

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Abstract: It is difficult to cut hard and brittle quartz glass chips. Hence, a method involving micro-crack-induced severing along a non-crack microgroove-apex by controlling the loading rate is proposed. The objective is to realize the rapid and precision severing of the hardest quartz glass in chip materials. Firstly, micro-grinding was employed to machine smooth microgrooves of 398–565 μm in depth; then the severing force was modelled by the microgroove shape and size; finally, the severing performance of a 4-mm thick substrate was investigated experimentally. It is shown that the crack propagation occurred at the same time from the microgroove-apex and the loading point during 0.5 ms in micro-crack-induced severing. The severing efficiency is dominated by the severing time rather than the crack propagation time. When the loading rate is less than 20–60 mm/min, the dynamic severing is transferred to static severing. With increasing microgroove-apex radius, the severing force decreases to the critical severing force of about 160–180 N in the static severing, but it increases to the critical severing force in the dynamic severing. The static severing force and time are about two times and about nine times larger than the dynamic ones, respectively, but the static severing form error of 16.3 μm/mm and surface roughness of 19.7 nm are less. It is confirmed that the ideal static severing forces are identical to the experimental results. As a result, the static severing is controllable for the accurate and smooth separation of quartz glass chips in 4 s and less.

Keywords: micro-crack propagation; severing force; quartz glass; micro-grinding

1. Introduction

Rolling scribing with a tungsten carbide or polycrystalline diamond (PCD) wheel is widely used to separate silicate glass substrates without any coolant and material removal. In order to improve the life of tool micro-tips, a chemical vapor deposition (CVD) diamond roller was developed for scribing instead of PCD and carbide alloy rollers [1]. Lateral and radial cracks were, however, produced on severing surfaces due to the mechanical force of mechanical wheel rolling [2,3]. Generally, a follow-on smooth profile grinding and polishing were needed along with inefficiency and pollution. Recently, researchers have focused on the crack generation on severing surfaces. From the in-process estimation of fracture surface morphology, severing surface cracks and breakages were produced during wheel scribing of a glass sheet [4]. It has been known that the median crack depth decreased with decreasing loading force, but it still reached about 90 μm with the load of 22 N in scribing alumina ceramics [1]. In scribing LCD glass, it reached about 45 μm with the load of 5.5 N [2] and about 90 μm with the cutting pressure of 0.16 MPa [3], respectively. Until now, the scribing of harder quartz glass chips has not yet been reported.

In order to improve scribing performance, vibration-assisted scribing was used to increase the median cracks for severing [5], but uneven cracks existed on microgroove-apex edges and the severing surface. It has been reported that laser beams could be used to irradiate the scribing [6–9]. Although
laser beam irradiation enhanced scribing speed, edge cracks, severing form deviation and the bevel surfaces were produced [6]. Moreover, the laser irradiation made the mechanical breaking easier in scribing [7], but thermal damage accumulated at the severing edges. Because the laser scribing produced micro-cracks and burrs on the machined microgrooves, it led to the cracks and burrs on the severing surface edge [8]. Although a hybrid of laser beam, water jet coolant and pre-bending were employed to eliminate the micro-cracks [9], form deviation happened along the beam moving direction. A picosecond Ultraviolet (UV) laser was also used to induce the scribing of polyethylene terephthalate films to control the local bending flexibility [10], but it has not yet been applied to severing.

To predict severing force, Filippi first proposed the existence of linear elastic stress fields in the neighborhood of rounded-tip V-shaped notches [11]. The linear elastic stress was also used to derive two brittle fracture criteria such as mean stress (MS) and point stress (PS) criteria [12]. Moreover, these criteria were used to predict compressive notch fracture toughness [13,14]. The minimum fracture loading of a U-notches plate was introduced by means of MS and PS criteria [15]. In the case of low and high loading rates, it was found that the loading rate produced little influence on the maximum load for the V-notch on fracture [16]. However, these workpieces only concerned easy-to-cut polymeric and metallic materials. Until now, these criteria have not been applied to difficult-to-cut quartz glass due to the fabrication difficulty of microgroove. Although the fracture of ceramic-metal joint surface has been divided into static and dynamic states [17], the critical loading rate and force have not yet been studied in detail.

As for the fabrication of microgroove, laser and etching approaches have been used to fabricate the microgroove with 7.5 µm and less in depth on Si surface and ceramic cylinder [18,19], but it was irregular and rough. It would lead to cracks on the severing surface when it was used for the induced severing. Moreover, the micro-grinding with a sharpened diamond wheel micro-tip may be employed to fabricate accurate and smooth microgrooves on difficult-to-cut silicon, carbide alloy and glass surfaces [20], but it has not yet been applied to the crack propagation for precision and smooth severing of difficult-to-cut materials.

In this paper, a new micro-crack-induced severing with static loading and dynamic loading is proposed for the crack propagation along an accurate and smooth microgroove-apex. The objective is to realize rapid and precision severing of difficult-to-cut quartz glass. Firstly, the trued diamond wheel micro-tip was employed to grind the accurate and non-crack microgroove on workpiece surface; then the severing force was modelled in micro-crack induced severing by microgroove parameters and loading rate; finally, severing force, severing time, cracking propagation time, severing form errors and severing surface roughness were experimentally investigated.

2. Micro-Crack Induced Severing of Brittle Workpiece

Figure 1 shows the stress field model in micro-crack induced severing along a microgroove-apex. The microgroove is parameterized by height $h$, angle $\beta$, and microgroove-apex radius $r_c$. Under mode I loading condition, $\sigma_{\theta\theta}$, $\sigma_{r\theta}$, and $\sigma_{rr}$ are tangential stress, shear stress and radial stress, respectively, and $r_c^*$ is critical distance (see Figure 1a) [11]. When the loading force $F$ increases to the critical value called severing force $F_c$, the tangential stress $\sigma_{\theta\theta}$ ($r_c^*$, 0) reaches the ultimate tensile strength $\sigma_p$, and the micro-cracks are produced from the microgroove-apex (see Figure 1b). It leads to the crack propagation along the microgroove-apex.

Figure 2 shows the scheme of micro-crack-induced severing. The working sizes of the workpiece substrate are given by the thickness $W$ and the width $B$. The workpiece is supported by two supporting rods with an interval $L$. The arc-shaped loading rod is loaded on the upper surface of substrate. The microgroove-apex is positioned on the opposite side of substrate. The vertical loading direction aims to the microgroove-apex. The loading rod moved vertically with the loading rate $v$. When the loading force $F$ reaches the severing force $F_c$, the micro-crack occurs at the microgroove-apex. It leads to the severing for the separation of workpiece.
3. Modelling of Severing Force

The elastic stresses were described at the neighborhood of microgroove-apex in polar coordinate system (see Figure 1a). \( r_0 \) is the distance between coordinate origin \( O \) and microgroove-apex. The critical distances \( r_c^* \) and \( r_c \) are defined from the coordinate origin \( O \) and the microgroove-apex, respectively (see Figure 1b). The ideal severing force \( F_c^* \) was described as follows [15]:

\[
\sigma_{\text{max}} = \sigma_{\text{nom}} \cdot K_f = \frac{3K_f F_c^* L}{2B(W - h_v)^2}
\]

where \( \sigma_{\text{max}} \) and \( \sigma_{\text{nom}} \) are the maximum stress and the nominal stress at microgroove-apex, respectively. \( K_f \) is the stress concentration factor.

At the neighborhood of microgroove-apex under pure mode I loading, the elastic stresses are described in the polar coordinate system as follows [11]:

\[
\begin{align*}
\sigma_\theta \theta &= \frac{K_f r v}{\sqrt{2\pi r^2}} \left[ m_{\theta \theta}(\theta) \right], \\
\sigma_r r &= \frac{K_f r v}{\sqrt{2\pi r^2}} \left[ m_{rr}(\theta) \right], \\
\sigma_\phi r &= \frac{K_f r v}{\sqrt{2\pi r^2}} \left[ m_{\rho \phi}(\theta) \right]
\end{align*}
\]

\[
\begin{bmatrix}
    m_{\theta \theta}(\theta) \\
    m_{rr}(\theta) \\
    m_{\rho \phi}(\theta)
\end{bmatrix} = \frac{1}{1 + \lambda_1 + \chi_2 (1 - \lambda_1)} \begin{bmatrix}
    (1 + \lambda_1) \cos[1 - \lambda_1] \theta \\
    (3 - \lambda_1) \cos[1 - \lambda_1] \theta \\
    (1 - \lambda_1) \sin(1 - \lambda_1) \theta
\end{bmatrix} + \lambda_2 (1 - \lambda_1) \begin{bmatrix}
    \cos(1 + \lambda_1) \theta \\
    -\cos(1 + \lambda_1) \theta \\
    \sin(1 + \lambda_1) \theta
\end{bmatrix}
\]

\[(3)\]
The $n_{th}(\theta)$ is expressed as:

$$
\begin{align*}
\begin{cases}
n_{th}(\theta) \\ n_{tr}(\theta) \\ n_{r}(\theta)
\end{cases} = \frac{1}{4(q-1)[1+\lambda_1+\chi_1(1-\lambda_1)]}
\begin{cases}
\lambda_1 (1+\mu_1) \cos(1-\mu_1) \theta \\ (3-\mu_1) \cos(1-\mu_1) \theta \\ (1-\mu_1) \sin(1-\mu_1) \theta
\end{cases} + \chi_1 (1+\mu_1) \theta
\end{align*}
$$

The $K_{I}^{V,r}$ is the mode I notch stress intensity factor (NSIF). It is described as follows:

$$
K_{I}^{V,r} = \frac{\sqrt{2\pi r_0 (1-\lambda_1)} \sigma_{max}}{1+\omega_1}
$$

where $\omega_1$ is an auxiliary parameter. They are expressed as follows:

$$
r_0 = \frac{q-1}{q} r_v
$$

$$
\omega_1 = \frac{q}{4(q-1)} \left[ \frac{\chi_1 (1+\mu_1) + \chi_1}{1+\lambda_1 + \chi_1 (1-\lambda_1)} \right]
$$

where $q$ is a real positive coefficient ranging as:

$$
q = \frac{2\pi - \beta_v}{\pi}
$$

$\lambda_1$, $\mu_1$, $\chi_1$, $\chi_1$ and $\chi_1$ are the values of auxiliary parameters for different microgroove angle [11].

Under pure mode I loading, the tangential stress $\sigma_{th}(r,0)$ from Equation (2) can be written as follows:

$$
\sigma_{th}(r,0) = \frac{K_{I}^{V,r}}{\sqrt{2\pi r_0 (1-\lambda_1)}} \left[ 1 + \left( \frac{r}{r_0} \right)^{\mu_1-\lambda_1} n_{th}(0) \right]
$$

Substituting Equation (5) into Equation (9), the tangential stress can be expression as follows:

$$
\sigma_{th}(r,0) = \frac{r_0^{1-\lambda_1} \sigma_{max}}{r^{1-\lambda_1} (1+\omega_1)} \left[ 1 + \left( \frac{r}{r_0} \right)^{\mu_1-\lambda_1} n_{th}(0) \right]
$$

According to PS criterion, the brittle fracture takes place when the tangential stress $\sigma_{th}(r,0)$ reaches critical value $\sigma_u$ at specified critical distance [21]. Hence, $r_c^*$ (see in Figure 1b) can be expressed as follows:

$$
r_c^* = r_c + r_0
$$

For brittle materials, the critical distance $r_c$ can be written as follows [22]:

$$
r_c = \frac{1}{8\pi} \left( \frac{K_{IC}}{\sigma_u} \right)^2
$$

where $K_{IC}$ is the material attribute called fracture toughness.

According to Equations (1) and (9)-(12), the ideal severing force $F_{c}^*$ is deduced as follows:

$$
F_{c}^* = \frac{2\sigma_u B(W-h_c)(1+\omega_1)(r_0 + r_c)^{1-\lambda_1}}{3K_L r_0^{1-\lambda_1} \left( 1 + \left( \frac{r_c}{r_0} \right)^{\mu_1-\lambda_1} n_{th}(0) \right)}
$$

In Equation (13), the $K_L$ is achieved in the case of the U-notch with $\beta_v = 0$ [15], but it was calculated by the fitting of experimental data in this study. This is because a microgroove with $\beta_v > 0$ was employed in micro-crack-induced severing.
4. Micro-Grinding of Microgroove on Workpiece Surface

Figure 3 shows the micro-grinding of microgrooves on a workpiece surface with a diamond wheel micro-tip. It was difficult to dress the wheel micro-tip due to its high hardness. Before micro-grinding, the truing of the #600 metal-bond diamond wheel micro-tip (Changxing technology Co. LTD., Shenzhen, China) was performed by the Numerical Control (NC) mutual wearing between the diamond wheel and the #800 Green Silicon Carbide (GC) dresser using Computer Numerical Control (CNC) grinder (Fuyu Machine Tool co. LTD., Zhang Hua, Taiwan) (see Figure 2a) [20]. Finally, the diamond wheel micro-tip angle $\alpha_v$ was identical to the angle of the NC tool paths. The truing conditions of the wheel micro-tip are shown in Table 1.

![Experimental scene](image)

Figure 3. Micro-grinding of microgroove on quartz glass: (a) the truing of diamond wheel micro-tip and (b) the micro-grinding of microgroove.

| CNC grinder | SMART B818 |
|-------------|------------|
| Diamond wheel | SD600 (Metal-bond, Grain size: 24 $\mu$m), $D = 150$ mm, $w = 4$ mm |
| Dresser | #800 GC, Ceramic bond |
| Tool paths | V-shaped symmetrical, $\alpha = 60^\circ$, $90^\circ$, $120^\circ$ |
| Truing parameters | $v_f = 500$ mm/min, $N = 2400$ rpm, $a = 20$ $\mu$m, $\sum a = 5$ mm |
| Coolant | Water |

Then, the diamond wheel micro-tip was used to grind a microgroove on the workpiece surface (see Figure 2b). A quartz glass substrate was chosen as workpiece. The diamond wheel micro-tip angle $\alpha_v$ was equal to the microgroove angle $\beta_v$ [19]. The micro-grinding conditions are shown in Table 2.
angle $\alpha_v$ was equal to the microgroove angle $\beta_v$ [19]. The micro-grinding conditions of microgroove on workpiece substrate.

Table 2.

| CNC Grinder | SMART B818 |
|-------------|-------------|
| Diamond wheel | SD600 (Metal-bond, grain size: 24 $\mu$m), $D = 150$ mm, $w = 4$ mm |
| Workpiece | Quartz glass |
| Grinding parameters | $v_f = 500$ mm/min, $N = 2400$ rpm, $a = 20$ $\mu$m, $\Sigma a = 500$ $\mu$m, |
| Coolant | Water |

5. Experiment and Measurement

Figure 4 shows the experimental setup of micro-crack induced severing. It is based on its working principle (see Figures 1 and 2). A WDW-05 electronic universal testing machine (Jinan Kason Testing Equipment Co., Ltd., Qingdao, Shandong, China) was employed to perform the loading and measure the on-line loading force $F$. A high-speed camera was used to record the propagation of micro-cracks during the loading process (see Figure 4). In the experiments, the microgroove angle $\beta_v$ was set as $60^\circ$, $90^\circ$ and $120^\circ$, respectively. The loading rate $v$ was set as $5$ mm/min, $10$ mm/min, $20$ mm/min, $60$ mm/min, $100$ mm/min, $200$ mm/min and $300$ mm/min, respectively. In order to calculate ideal severing force $F_c^*$, the values of the auxiliary parameters such as $\chi_1$, $\mu_1$, $\chi_{c1}$ and $\lambda_{d1}$ in Equation (13) were given in Table 3. The mechanical properties of quartz glass were given in Table 4.

![Figure 4](image_url)

Figure 4. The experimental setup of micro-crack induced severing.

Table 3. The auxiliary parameters.

| $\beta_v$ | $\lambda_1$ | $\mu_1$ | $\chi_{b1}$ | $\chi_{c1}$ | $\lambda_{d1}$ |
|----------|-------------|---------|-------------|-------------|-------------|
| $60^\circ$ | 0.5122 | $-0.4057$ | 1.3123 | 3.2832 | 0.0960 |
| $90^\circ$ | 0.5448 | $-0.3449$ | 1.8414 | 2.5057 | 0.1046 |
| $120^\circ$ | 0.6157 | $-0.2678$ | 3.0027 | 1.5150 | 0.0871 |

Table 4. Mechanical properties of quartz glass.

| Property | Value |
|----------|-------|
| Young’s modulus [GPa] | 77 |
| Poisson’s ratio | 0.17 |
| Ultimate tensile strength $\sigma_u$ [MPa] | 50 |
| Fracture toughness $K_{IC}$ [MPa-m$^{0.5}$] | 0.81 |

In order to compare the traditional scribing with the micro-crack induced severing, an SFT-QG500A glass-cutter machine (Shufeng technology co. Ltd., Shenzhen, China) was employed.
to perform the scribing-and-breaking severing experiments of quartz glass. The rolling scribing conditions were given by 120° in tungsten carbide wheel V-tip angle, 0.19 MPa in scribing pressure, 1 mm in setting depth and 30 mm/s in scribing speed.

6. Results and Discussions

6.1. Profile of Micro-Ground Microgroove on Workpiece Surface

Figure 5 shows the microgroove profile of quartz glass after micro-grinding. It is shown that the microgroove was regular and its edges were smooth. No micro-cracks existed on the microgroove. This is because the grain cutting depths may be controlled to be less than the critical cutting depth transferred from brittle cutting to ductile cutting in micro-grinding, leading to a no-crack microgroove. According to the measured results of VHX-1000 microscope (Keyence, Osaka, Japan), the microgroove angle $\beta_v$ averagely reached 61.7°, 91.8° and 119.4° in contrast to the designed ones of 60°, 90° and 120°, respectively. The microgroove angle error was ±1.8°. Correspondingly, the microgroove heights $h_v$ were 545 μm, 495 μm and 479 μm, respectively. The microgroove-apex radius $r_v$ were 39.3 μm, 41.3 μm and 39.7 μm, respectively. The microgroove surface roughness $R_a$ was 80–100 nm. As a result, the micro-grinding was able to fabricate an accurate and smooth microgroove without any cracks around its apex. In contrast, the scribing produced the cracks along the scratch [2,3].

![Figure 5. The microgroove profile of quartz glass after micro-grinding.](image)

6.2. Severing Surface Topography

Figure 6 shows the severing surface topographies of quartz glass in mechanical rolling scribing and micro-crack-induced severing. It is shown that breakages happened on the scratch microgroove edges in mechanical rolling scribing (see Figure 6a). This is because the pure mechanical compression produced the radial cracks, leading to edge cracks. It also produced median cracks, radial cracks and lateral cracks, leading to an uneven severing surface in break-severing. It was identical to the results in the scribing-and-breaking severing of ceramics and silica glass substrates [1,3]. In contrast, no breakages happened on the severing workpiece edges along the micro-ground microgroove-apex in micro-crack-induced severing (see Figure 6b).

The severing surface edges were undamaged. The severing surface was flat and smooth. This is because the crack propagation was precisely induced along the micro-ground microgroove-apex in micro-crack-induced severing (see Figure 5). This also means that the accurate and smooth microgroove-apex without any micro-cracks was able to induce the accurate and smooth crack propagation in severing.
Figure 6. The severing surface topographies of quartz glass substrate: (a) mechanical rolling scribe (b) micro-crack induced severing.

6.3. Loading Force and Loading Time versus Loading Rate

Figure 7 shows the loading force $F$ and loading time $T$ versus loading rate $v$. Experimental results showed that the loading force $F$ increased with increasing loading time $T$ at beginning, but it rapidly decreased after the micro-crack propagation happened on the microgroove-apex.

It was identical to the relationship between loading force and displacement in tensile fracture of notched polycrystalline graphite [23]. The critical loading force and time were regarded as severing force $F_c$ and severing time $T_c$, respectively. In the case of loading rate $v = 60–300$ mm/min, the severing force rapidly increased with increasing loading time $T$, but it slowly increased in the case of $v = 5–20$ mm/min. This mean that there existed two different mechanisms in micro-crack induced
severing. Their critical loading rate \( v_0 \) was distributed between 20 mm/min and 60 mm/min. As a result, the micro-crack induced severing may be distinguished by static severing \( (v < v_0) \) and dynamic severing \( (v > v_0) \), respectively.

6.4. Severing Force versus Microgroove-Apex Radius and Loading Rate

Figure 8 shows the severing force \( F_c \) versus microgroove-apex radius \( r_\alpha \) for different loading rate \( v \) and microgroove angle \( \beta_v \). Three experiments were accomplished for the same loading rate. The singularity was removed. It is shown that the severing force \( F_c \) rapidly increased and slowly approach the critical severing force \( F_{cc} \) with increasing microgroove-apex radius \( r_\alpha \) in dynamic severing, but it slowly decreased and gradually approach the critical severing force \( F_{cc} \) in static severing. This mean that the static severing and dynamic severing produced different influence on severing force with reference to microgroove-apex radius. Moreover, the microgroove angle \( \beta_v \) produced little influence on the severing force \( F_c \). The severing force \( F_c \) averagely reached 226.5 N in static severing, but it averagely reached 116.2 N in dynamic severing. As a result, the static severing increased the severing force by 95% compared to the dynamic severing.

![Figure 8. Severing force \( F_c \) versus microgroove-apex radius \( r_\alpha \) for different loading rate \( v \) and microgroove angle \( \beta_v \).](image)

It is also seen that the severing force \( F_c \) rapidly decreased with increasing loading rate \( v \) in dynamic severing. In contrast, it slowly decreased in static severing. However, the loading rate produced little influence on maximum loading force in low and high loading rate when the workpiece was thermoset epoxy resin [14]. When the microgroove-apex radius was larger than the critical value of 40–50 \( \mu \)m, the severing force was not dominated by the microgroove-apex radius and the loading rate. Moreover, the critical severing force \( F_{cc} \) ranged 160–180 N when the dynamic severing was transformed into static severing.

6.5. Prediction of Severing Force

In contrast to the ideal \( F_{c*} \), the experimental severing force \( F_c \) was fitted to achieve \( K_t \) as follows:

\[
K_t = -0.3741 \beta_v + 4.787
\]  

(14)

According to Equation (13), the ideal severing force \( F_{c*} \) may be described as follows:

\[
F_{c*} = \frac{326.06 (0.4012r_\alpha + 0.0104)^{0.4878}}{r_\alpha^{0.4878} \left( 1 + 0.5791 \left( 1 + \frac{0.0104}{0.4012r_\alpha} \right)^{-0.9179} \right)^{0.4878}} \beta_v = 60^\circ
\]

(15)
When the microgroove-apex radius \( r \) was larger than the critical value of 40–50 \( \mu \text{m} \) (see Figure 8), the results of ideal severing force \( F_c^* \) were plotted in Figure 8. It is seen that the experimental cutting force was identical to the ideal severing force \( F_c^* \) in static severing, but it produced little influence on severing time in dynamic severing. Moreover, the severing force \( F_c^* \) was stabilized at 159.5–190.0 N, which was identical to the experimental results. This mean that the Equations (13)–(16) may be used to predict and control the severing force according to the microgroove height \( h_v \), angle \( \beta_v \) and microgroove-apex radius \( r_v \) in static severing.

### 6.6. Ideal Severing Force versus Microgroove Angle and Height

Figure 9 shows the ideal severing force \( F_c^* \) versus microgroove angle \( \beta_v \) and height \( h_v \) in static severing. It is shown that the microgroove angle \( \beta_v \) produced little influence on the severing force \( F_c^* \) in the case of \( \beta_v = 60–120^\circ \), but the severing force \( F_c^* \) decreased with increasing microgroove \( \beta_v \) in the case of \( \beta_v > 120^\circ \) and \( h_v > 30 \mu \text{m} \), respectively (see Figure 9a). It is also seen that the severing force \( F_c^* \) decreased with increasing microgroove height \( h_v \) and microgroove-apex radius \( r_v \) (see in Figure 9b). When the microgroove-apex radius \( r_v \) was larger than 40–50 \( \mu \text{m} \), the severing force was not dominated by the microgroove angle and microgroove-apex radius.

\[
F_c^* = \frac{331.66(0.3333 \beta_v + 0.0104)^{0.4552}}{r_v^{0.4552} \left( 1 + 0.54 \left( 1 + \frac{0.0104}{0.3333 \beta_v} \right)^{-0.8897} \right)}, \beta_v = 90^\circ
\]

\[
F_c^* = \frac{313.83(0.2481 \beta_v + 0.0104)^{0.3843}}{r_v^{0.3843} \left( 1 + 0.4319 \left( 1 + \frac{0.0104}{0.2481 \beta_v} \right)^{-0.8835} \right)}, \beta_v = 120^\circ
\]

Using Equations (15)–(17), the results of ideal severing force \( F_c^* \) were plotted in Figure 8. When the microgroove-apex radius \( r \) was larger than the critical value of 40–50 \( \mu \text{m} \), the ideal severing force \( F_c^* \) was stabilized at 159.5–190.0 N, which was identical to the experimental results. This mean that the Equations (13)–(16) may be used to predict and control the severing force according to the microgroove height \( h_v \), angle \( \beta_v \) and microgroove-apex radius \( r_v \) in static severing.

### 6.7. Severing Time versus Microgroove-Apex Radius

Figure 10 shows the severing time \( T_c \) versus microgroove-apex radius \( r_v \). The severing time \( T_c \) was regarded as the mean value of experimental data at the same loading rate \( v \). Experimental results showed that the severing time \( T_c \) rapidly decreased with increasing microgroove-apex radius \( r_v \) in static severing, but the microgroove-apex radius \( r_v \) produced little influence on severing time in dynamic severing. Moreover, the severing time \( T_c \) averagely reached 2.43 s in static severing, but it averagely reached 0.27 s in dynamic severing. Hence, the static severing increased the severing time by about 900% compared to the dynamic severing.
Figure 10. Severing time $T_c$ versus microgroove-apex radius $r_c$.

6.8. Severing form Errors versus Loading Rate

Figure 11 shows the microgroove-direction $e_m$ and the loading-direction severing form errors $e_l$ versus loading rate $v$ in micro-crack induced severing. Experimental results showed that the severing form errors gradually increased with increasing loading rate $v$ in both static severing and dynamic severing. The microgroove-direction severing form error $e_m$ averagely reached 8.8 μm/mm (see Figure 11a), which was much less than the loading-direction severing form error of 31.7 μm/mm (see Figure 11b). This is because the microgroove direction was dominated by the accurate micro-ground microgroove-apex, but the loading direction depended on precision positioning and loading rate. Moreover, the static severing form error averagely reached 16.2 μm/mm, which was less than the dynamic severing form error of 25.3 μm/mm. The reason is that high loading rate easily produced the position deviation between workpiece and loading rod. Hence, the static severing may decrease the severing form error by 36% compared to the dynamic severing.

6.9. Severing Surface Roughness versus Loading Rate

Figure 12 shows the severing surface roughness $R_s$ versus loading rate $v$. Experimental results showed that the loading rate $v$ had little influence on the severing surface roughness $R_s$ in both static severing and dynamic severing. However, the microgroove-direction severing surface roughness of 13.7 nm was much less than the loading-direction severing surface roughness of 29.6 nm. This means that severing surface roughness could be dominated by material properties. Moreover, the static severing surface roughness of 19.69 nm was less than the dynamic severing surface roughness of 22.34 nm. Hence, the static severing surface roughness decreased by 12% compared to the dynamic...
severing. As a result, the micro-crack induced severing may produce smooth severing surface of quartz glass without any polishing.

Figure 12. Severing surface roughness $R_a$ versus loading rate $v$: (a) microgroove-direction severing roughness and (b) loading-direction severing roughness.

6.10. Cracking Propagation Time and Loading Positions

Figure 13 shows the micro-crack propagation process in static severing. It is shown that crack propagation occurred at the same time from the microgroove-apex and the loading point in 0.5 ms (see Figure 13a,b). This mean that the micro-crack was induced from the microgroove-apex to extend to the loading point. Hence, the loading positions dominated the severing form errors in micro-crack induced severing. This may explain why the severing form errors increased with increasing loading rate $v$ (see Figure 11).

Figure 13. Micro-crack propagation process in static severing ($v = 20$ mm/min).

It is also found that the cracking propagation time reached 0.5 ms and less to achieve the smooth surface (see Figure 13b). In contrast, the severing time $T_c$ averagely reached 2430 ms in static severing (see Figure 10). This also mean that the severing time was much larger than the crack propagation time. Hence, the efficiency of micro-crack induced severing is dominated by the severing time rather than the crack propagation time.

7. Conclusions

1. Compared to a mechanical rolling scribe, the micro-crack-induced severing by a non-cracked microgroove-apex produces smooth severing edges without any breakages, median cracks,
radial cracks and lateral cracks. The micro-grinding may machine the accurate and smooth microgroove-apex without any cracks.

2. In micro-crack-induced severing, the crack propagation occurred at the same time from the microgroove-apex and the loading point. The severing efficiency is dominated by the severing time rather than the crack propagation time. The severing energy and quality depend on the loading rate.

3. When the loading rate was less than 20–60 mm/min, the dynamic severing changes to static severing. In static severing, the severing force slowly decreases with increasing microgroove-apex radius, but it rapidly increases in dynamic severing. For the critical severing force of 160–180 N, the static severing force is about two times larger than the dynamic one.

4. The severing time rapidly decreases with increasing microgroove-apex radius in static severing, but it slowly increases in dynamic severing. It reaches on average 2.43 s and 0.27 s, respectively. The static severing increases the severing time by about 9 times compared to the dynamic severing.

5. The severing form error and the severing surface roughness reach 4.6–32.5 µm/mm and 12.8–34.9 nm in static severing and 6.9–56.3 µm/mm and 13.3–40.9 nm in dynamic severing, respectively. The static severing may decrease the severing form error by 36% and the severing surface roughness by 12% compared to the dynamic severing, respectively.

6. In static severing, the severing force may be modelled and predicted by microgroove-apex radius, microgroove angle and height. Theoretically, it decreases with increasing the microgroove height. In dynamic severing, it is little influenced by microgroove-apex radius and microgroove angle when the microgroove angle is less than 120°.

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