EVIDENCE FOR EVOLUTION AMONG PRIMORDIAL DISKS IN THE 5 Myr OLD UPPER SCORPIUS OB ASSOCIATION

S. E. Dahm
W. M. Keck Observatory, 65-1120 Mamalahoa Highway, Kamuela, HI 96743, USA
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ABSTRACT

Moderate-resolution, near-infrared spectra between 0.8 and 5.2 μm were obtained for 12 late-type (K0–M3) disk-bearing members of the ~5 Myr old Upper Scorpius OB association using SpeX on the NASA Infrared Telescope Facility. For most sources, continuum excess emission first becomes apparent between ~2.2 and 4.5 μm and is consistent with that produced by single-temperature blackbodies having characteristic temperatures ranging from ~500 to 1300 K. The near-infrared spectra for 5 of 12 Upper Scorpius sources exhibit Paβ, Paγ, and Brγ emission, indicators of disk accretion. Using a correlation between Paβ and Brγ emission line luminosity and accretion luminosity, mass accretion rates (\( \dot{M} \)) are derived for these sources that range from \( M = 3.5 \times 10^{-10} \) to \( 1.5 \times 10^{-8} \, M_\odot \, yr^{-1} \). Merging the SpeX observations with Spitzer Space Telescope mid-infrared (5.4–37.0 μm) spectroscopy and 24 and 70 μm broadband photometry, the observed spectral energy distributions (SEDs) are compared with those predicted by two-dimensional, radiative transfer accretion disk models. Of the nine Upper Scorpius sources examined in this analysis, three exhibit SEDs that are most consistent with models having inner disk radii that substantially exceed their respective dust sublimation radii. The remaining Upper Scorpius members possess SEDs that either show significant dispersion among predicted inner disk radii or are best described by models having inner disk rims coincident with the dust sublimation radius.

Key words: accretion, accretion disks – open clusters and associations: individual (Upper Scorpius OB Association) – stars: formation – stars: pre-main sequence

Online-only material: color figures

1. INTRODUCTION

A fundamental consequence of protostellar collapse is the formation of a viscous accretion disk that transfers gas onto the stellar photosphere. Such disks are the progenitors of planetary systems and their subsequent evolution and ultimate dispersal have extraordinary implications for the formation of planets, orbital dynamics of planetary mass bodies, and the processing of dust grains and volatiles within the disk. The timescale of disk dissipation within the terrestrial region has been reasonably well established by ground-based (Haisch et al. 2001; Mamajek et al. 2004) and Spitzer Space Telescope (Uchida et al. 2004; Silverstone et al. 2006) observations to be \( \lesssim 10 \) Myr. By inference, planetary systems must be in an advanced evolutionary stage within this period, before significant quantities of gas and dust are depleted from the disk. Disk evolution is expected to proceed from the interior outward (Dullemond & Dominik 2005), however, submillimeter and mid-infrared observations by Cieza et al. (2008) suggest that inner disks only begin to dissipate after the outer disk has been significantly depleted of mass.

Pre-main-sequence stars exhibiting spectral energy distributions (SEDs) that are indicative of an optically thin disk interior surrounded by an optically thick outer disk are referred to as transition disk objects (Strom et al. 1989; Najita et al. 2007; Muzerolle et al. 2010). Such disks are suggestive of having experienced significant evolution from the continuous disk structures that are associated with classical T Tauri stars (CTTSs). The duration of this transitional phase has been inferred from population statistics of star-forming regions to be only of order \(~10^5\) yr (Hartmann 2009; Luhman et al. 2010). Transition-like SEDs, however, can arise from various pathways and may not be representative of a direct evolutionary sequence from primordial disk to debris disk (Najita et al. 2007). Dust grain growth and mid-plane settling (Dullemond & Dominik 2005), giant planet formation and dynamical clearing (Calvet et al. 2002), disk photoevaporation (Alexander et al. 2006), and the presence of a stellar companion (Ireland & Kraus 2008) have been suggested as disk clearing mechanisms capable of producing transition-like SEDs. An alternate evolutionary path from primordial disk to debris disk has been proposed by Currie et al. (2009), who suggest that disks exhibiting reduced levels of near- and mid-infrared excess emission are indicative of reduced masses of small dust grains at all disk radii. Such homologously depleted disks counter the canonical inside-out disk evolutionary scenario.

The Upper Scorpius OB association is a critically important region for studies of disk evolution. At ~145 pc distant, it is among the nearest OB associations to the Sun (Blauuw 1991; de Zeeuw et al. 1999) and has a well-established age of ~5 Myr (Preibisch & Zinnecker 1999; Preibisch et al. 2002), when most (~80%) optically thick, protoplanetary disks have dissipated (Haisch et al. 2001; Hernandez et al. 2007). Carpenter et al. (2006) conducted a Spitzer 4.5–16.0 μm photometric survey of 218 confirmed association members. These sources were compiled from Hipparcos astrometry (de Zeeuw et al. 1999), color–magnitude diagrams and Li i \( \lambda \) 6708 follow-up observations (Preibisch & Zinnecker 1999; Preibisch et al.
The 12 Upper Scorpius infrared excess sources observed with SpeX were selected from the mid-infrared photometric and spectroscopic surveys of Carpenter et al. (2006) and Dahm & Carpenter (2009), respectively. These sources represent half of the late-type, disk-bearing stars identified by Carpenter et al. (2006) and include all K-type and all M0 through M2-type excess sources. To further place this sample into context with the low-mass stellar population of Upper Scorpius, there are ~250 known pre-main-sequence stars in the mass range from ~0.1 to ~2.0 $M_\odot$ (Preibisch & Mamajek 2008; Preibisch et al. 2002). If the infrared excess fraction for late-type stars (19%) derived by Carpenter et al. (2006) is assumed, ~48 of these should host primordial disks. The SpeX sample therefore represents ~25% of all disk-bearing stars expected among the known late-type association members. Integrating the best-fitting mass function for Upper Scorpius, Preibisch & Mamajek (2008) estimate that the total low-mass (0.1−2.0 $M_\odot$) stellar population of the association exceeds 2400 stars, corresponding to a total disk-bearing population of ~450 late-type stars. While the SpeX sample considered here does represent a substantial fraction of the late-type, disk-bearing stars identified by Carpenter et al. (2006), it cannot be considered a statistically significant representation of the total low-mass stellar population of the association.

Given that unresolved companions could provide an explanation for transition-like SEDs (Ireland & Kraus 2008), some knowledge of the binary frequency of the Upper Scorpius late-type stellar population is warranted. Köhler et al. (2000) used speckle interferometry and direct imaging for 118 X-ray selected T Tauri stars in the greater Scorpius-Centaurus OB association to identify companions with separations ranging from 0′′13 to 6′′0. They found a multiplicity fraction of ~32.6%, which exceeds the binary fraction of main-sequence stars by a factor of ~1.6, but is slightly lower than that observed in Taurus-Auriga (Köhler et al. 2000). Kraus et al. (2008) found a similar frequency of binary companions among 82 late-type (G0–M4) Upper Scorpius members, ~35% ± 5%. Two of the disk candidates observed with SpeX are established binaries, [PZ99]J161411.1-230536 and ScoPMS 31, resolved by high angular resolution imaging. With projected separations of ~32 and 84 AU, respectively (Metchev & Hillenbrand 2009; Köhler et al. 2000), it is unlikely that these stellar companions have significantly influenced the inner disk evolutionary timescales of their primaries. Kraus et al. (2008) observed J160823.2-193001 and J160900.7-190852 using an aperture mask interferometry technique, placing firm upper limits that preclude the possibility of an undetected stellar companion within ~5 AU of these stars. The multiplicity fraction among the remaining Upper Scorpius disk-bearing sample is relatively unexplored, but is not expected to deviate significantly from that derived by Köhler et al. (2000) and Kraus et al. (2008) for the pre-main-sequence population as a whole.

Shown in Figure 1 is the $J - H$, $[8.0] - [4.5]$ color–color diagram for the 218 Upper Scorpius members included in the Carpenter et al. (2006) Spitzer IRAC and IRS survey. The 12 sources observed with SpeX span the full range of the 8.0 $\mu m$ excess distribution, from moderate excess (e.g., J160643.8-190805, ScoPMS 31) to significant (e.g., J160900.7-190852). Their masses, derived using the pre-main-sequence models of Siess et al. (2000) and assuming a distance of 145 pc, range from ~0.25 to 2.0 $M_\odot$. General properties of the stellar
sample including spectral type, extinction ($A_V$), mass, radius, luminosity, and effective temperature are presented in Table 1.

2.2. SpeX Observations

The SpeX (Rayner et al. 2003) observations were made on the nights of 2009 May 20–23 under variable cirrus and seeing conditions (0.6–0.9). The Upper Scorpius members were observed in the short cross-dispersed (SXD, 0.84–2.4 μm) and long cross-dispersed (LXD, 2.2–5.2 μm) modes for complete near-infrared wavelength coverage. Bright A0 V stars were observed at similar airmasses ($Δμ < 0.1$) and within short periods of time of the program stars for telluric correction. Arc lamps and internal flat-field exposures were obtained for each set of observations to account for instrument flexure. All observations were made with the 0.5 slit yielding a nominal spectral resolution of $λ/Δλ ≈ 1500$.

The spectra were reduced using SpeXtool, an IDL-based reduction package that provides for sky subtraction, flat-fielding, wavelength calibration, and optimal extraction (Cushing et al. 2004). For telluric corrections $xtellcor$, an extension package of SpeXtool, was used to create a kernel for the telluric spectrum using a model spectrum of α Lyra. $xtellcor$ interpolates over broad hydrogen absorption lines in the spectra of the A0 V stars using a technique developed by Vacca et al. (2003). The orders of the telluric corrected spectra were then combined and the SXD and LXD spectra merged using routines available within SpeXtool. The resulting continuum levels across all orders and between the SXD and LXD modes were found to be remarkably consistent. The near-infrared spectra were then merged with the Spitzer IRS spectra. Details of the Spitzer IRS observations and their subsequent reduction and analysis can be found in Dahm & Carpenter (2009).

3. NEAR-INFRARED CONTINUUM EXCESS SPECTRA

The extinction-corrected, near-infrared (0.84–5.2 μm) spectra of the 12 Upper Scorpius sources are shown in Figure 2. Regions of strong telluric absorption (atmospheric transmission <20%) have been excised from the figure. Also shown are the spectra of solar metallicity, main-sequence stars of identical or closely matched ($±1$ sub-class) spectral type obtained from the IRTF spectral library (Rayner et al. 2009). These standards have reliable spectral types and are directly linked to the Morgan–Keenan classification system. The standard or template spectra are scaled to the flux levels of the Upper Scorpius sources at 1.65 μm, near the peak of the stellar SED and where extinction effects are minimized (D’Alessio et al. 1999; Furlan et al. 2006).

To characterize excess emission attributable to the inner disk rim, the scaled photospheric template spectra are subtracted from the reddened spectra of the Upper Scorpius sample. The resulting continuum excess spectra are shown in Figure 3. Significant noise is present due to imperfect telluric correction, particularly in the thermal region where the water column varies both temporally and with airmass. In general, the slopes of the continuum excess spectra are increasing from $K$ band to $≈3.5$ μm before turning over toward redder wavelengths. The shapes of many of the continuum excess spectra are suggestive of having been produced by single-temperature blackbodies. By fitting these excess spectra with Planck functions, the characteristic temperatures of the blackbodies are found to range from the sublimation temperature for silicate dust, i.e., $≈1400$ K, to $<500$ K. The slope of excess emission, particularly between $≈2.75$ and 4.2 μm, was found to be of critical importance when fitting the blackbody profiles to the continuum excess spectra.

To estimate the uncertainty associated with the derived blackbody temperatures, photospheric templates within a range of $±1$ sub-class of the assigned spectral type were subtracted from the Upper Scorpius spectra. The best-fitting blackbody curves for this range of photosphere-subtracted excess spectra suggest typical uncertainties of $≈±200$ K for most sources. The adopted blackbody curve as well as curves representing the limits of uncertainty in the characteristic temperature are shown in Figure 3. The characteristic temperatures of the best-fitting blackbodies are provided in Table 1 and are assumed to be representative of the dust temperatures in the disk interiors.

The disk fractions of Upper Scorpius and Taurus-Auriga differ significantly for the range in stellar masses considered here: 19% for the former (Carpenter et al. 2006) and up to 75% for the latter (Luhman et al. 2010). Muzerolle et al. (2003) used SpeX in LXD mode to observe nine CTTSs in Taurus-Auriga, eight of which span a range of spectral types (G5–M1), luminosities (12.8–0.5 $L_⊙$), and stellar radii (3.6–1.8 $R_⊙$) that are comparable with those of the Upper Scorpius disk sample. Restricting further comparison of these samples to just those sources having similar physical properties (i.e., spectral type, luminosity), the spectral profiles of continuum excess emission of the Upper Scorpius disks are found to differ from those of Class II sources in Taurus-Auriga. For six of nine members in the reduced Upper Scorpius sample, the characteristic temperatures of excess emission fall below the dust sublimation
temperature, a condition satisfied by only one member (the binary DQ Tau) of the Taurus-Auriga sample. Furthermore, the excess profiles of J160643.8-190805 and J160357.9-194210, which are not represented in Figure 3, exhibit no distinctive shape, implying minimal continuum excess emission at near-infrared wavelengths. Shallow rises may be present near ∼3.75 μm for these two sources, suggestive of warm (≤500 K) dust emission, but higher supernova (SN) spectra are needed
Figure 3. Near-infrared continuum excess spectra for 10 Upper Scorpius late-type, disk-bearing stars created by subtracting scaled photospheric template spectra obtained from the IRTF spectral library (Rayner et al. 2009) from the dereddened object spectra. The template spectra are of identical spectral type of the Upper Scorpius sources or closely matched, within ±1 sub-class. Significant noise is present due to imperfect telluric correction, particularly in the thermal region where the water column varies both temporally and with airmass. Shown in red are the Planck functions that best fit the continuum excess spectra. Most have characteristic temperatures ranging from the dust sublimation temperature, near \( \sim 1400 \) K, to less than \( \sim 500 \) K. Also depicted (in blue) are blackbody curves representing the upper and lower limits of uncertainty associated with the blackbody characteristic temperatures.

(A color version of this figure is available in the online journal.)

for confirmation. The excess spectra of J161420.2-190648, [PZ99]J160357.6-203105, J160900.7-190852, and J155829.8-231007 all have characteristic temperatures at or near the dust sublimation temperature, implying that emission arises from dust in close proximity to the host star. Of these sources, all are suspected accretors (Section 4).

4. ACCRETION LUMINOSITY AND MASS ACCRETION RATES

Dahm & Carpenter (2009) identified five Upper Scorpius members in the present sample as probable accretors using the H\( \alpha \) velocity width criteria of White & Basri (2003): [PZ99]J160357.6-203105 (K5), J160900.7-190852 (K7), J161420.2-190648 (M0), ScoPMS 31 (M0.5), and J155829.8-231007 (M3). The near-infrared spectra of these sources exhibit He \( \text{i} \) \( \lambda \)10830, Pa\( \gamma \), and Pa\( \beta \) emission and for most, weak Br\( \gamma \) emission. Shown in Figure 4 are the normalized spectra of these sources centered upon Pa\( \gamma \), Pa\( \beta \), and Br\( \gamma \). Dahm & Carpenter (2009) used veiling at \( \lambda 6500 \) Å as well as Ca \( \text{ii} \) \( \lambda 8542 \) emission line luminosity to estimate mass accretion rates (\( \dot{M} \)) for these stars. In the near-infrared, Muzerolle et al. (1998) found Pa\( \beta \) and Br\( \gamma \) emission line luminosity (\( L_{\text{Pa\( \beta \)}} \) and \( L_{\text{Br\( \gamma \)}} \)) to be well correlated with accretion luminosity, \( L_{\text{acc}} \). The resulting least-squares fits from Muzerolle et al. (1998) are given by

\[
\log \left( \frac{L_{\text{acc}}}{L_{\odot}} \right) = (1.14 \pm 0.16) \log \left( \frac{L_{\text{Pa\( \beta \)}}}{L_{\odot}} \right) + 3.15 \pm 0.58 \tag{1}
\]

\[
\log \left( \frac{L_{\text{acc}}}{L_{\odot}} \right) = (1.14 \pm 0.16) \log \left( \frac{L_{\text{Pa\( \beta \)}}}{L_{\odot}} \right) + 3.15 \pm 0.58 \tag{1}
\]
Table 1

| Source          | SpT  | $A_V$ (mag) | Mass ($M_\odot$) | Radius ($R_\odot$) | Luminosity ($L_\odot$) | $T_{\text{eff}}$ (K) | $T_D$ (K) |
|-----------------|------|-------------|------------------|---------------------|------------------------|----------------------|-----------|
| [PZ99]J161411.0-230536 | K0   | 2.4         | 1.98             | 2.65                | 5.62                   | 5329                 | 700       |
| [PZ99]J160421.7-213028  | K2   | 1.0         | 1.08             | 1.14                | 0.74                   | 4893                 | 900       |
| [PZ99]J160357.6-203105   | K5   | 0.9         | 1.09             | 1.59                | 0.82                   | 4321                 | 1300      |
| J160643.8-190805         | K6   | 1.9         | 0.95             | 1.42                | 0.65                   | 4205 $\leq$500      |           |
| J160823.2-193001         | K9   | 1.5         | 0.70             | 1.23                | 0.37                   | 3965                 | 800       |
| J160900.7-190852         | K9   | 0.8         | 0.69             | 1.24                | 0.43                   | 3965                 | 1300      |
| J161420.2-190648         | M0   | 1.8         | 0.56             | 1.55                | 0.52                   | 3840                 | 1300      |
| ScoPMS 31               | M0.5 | 0.9         | 0.52             | 1.63                | 0.60                   | 3782                 | 900       |
| J161115.3-175721         | M1   | 1.6         | 0.48             | 1.55                | 0.48                   | 3725                 | 900       |
| J160357.9-194210         | M2   | 1.7         | 0.40             | 1.06                | 0.17                   | 3611 $\leq$500      |           |
| J160545.4-202308         | M2   | 2.2         | 0.36             | 1.05                | 0.16                   | 3530                 | 800       |
| J155829.8-231007         | M3   | 1.3         | 0.25             | 0.60                | 0.05                   | 3380                 | 1300      |

Notes.

a Identifiers from Carpenter et al. (2006).

b Spectral type from the literature.

c Extinction estimates are taken from Preibisch & Zinnecker (1999) and Preibisch et al. (2002).

d From the models of Siess et al. (2000), assuming a distance of 145 pc.
e Dust temperature derived from the blackbody fits of the continuum excess emission.

and

$$\log \left( \frac{L_{\text{acc}}}{L_\odot} \right) = (1.26 \pm 0.19) \log \left( \frac{L_{\text{Br} \gamma}}{L_\odot} \right) + 4.43 \pm 0.79. \quad (2)$$

For the five suspected accretors in the Upper Scorpius sample, template spectra of identical or similar spectral type from the IRTF spectral library of Rayner et al. (2009) were used to subtract Paβ and Brγ photospheric absorption. Paβ and Brγ line luminosities were then determined using their measured equivalent widths and the extinction-corrected J- and Ks-band magnitudes obtained from the 2MASS point source catalog. The resulting emission line luminosities were then transformed into $L_{\text{acc}}$ using the above linear relationships of Muzerolle et al. (1998). The derived $L_{\text{acc}}$ values are directly proportional to $\dot{M}$ such that

$$L_{\text{acc}} \sim \frac{GM_\star \dot{M}}{R_\star} \left( 1 - \frac{R_\star}{R_{\text{in}}} \right), \quad (3)$$

where the stellar mass and radius estimates used are those predicted by the pre-main-sequence models of Siess et al. (2000). The factor of $(1 - R_\star/R_{\text{in}})$ is assigned a value of 0.8, which assumes an inner disk radius ($R_{\text{in}}$) of 5 $R_\star$ (Gullbring et al. 1998). Given the possibility that inner disk radii for the Upper Scorpius sources may exceed those of typical Class II sources in Taurus-Auriga, this value could be underestimated by a factor of $\sim$1.25. The derived $L_{\text{acc}}$ and $\dot{M}$ values with their associated uncertainties are presented in Table 2. These uncertainties arise from multiple sources: error in the measured equivalent width (assumed to be $\sim$20%), spectral type uncertainty when correcting for photospheric absorption ($\pm$1 sub-class), the uncertainty in each coefficient of the Muzerolle et al. (1998) relationships, and the uncertainty in $R_{\text{in}}$ when determining $\dot{M}$ values.

In general the $\dot{M}$ values derived using $L_{\text{Paβ}}$ and $L_{\text{Brγ}}$ agree reasonably well with each other and with the $\dot{M}$ values from the veiling and Ca II λ8542 analysis of Dahm & Carpenter (2009). In summary, 5 of 12 of the Upper Scorpius sample are accreting, providing unambiguous evidence for the presence of gas within the terrestrial regions of these disk-bearing systems.

Figure 4. Emission features of He i λ10830 and Paγ (left), Paβ (center), and Brγ (right), for the five accreting sources in the sample. From top to bottom, the sources shown are ScoPMS 31 (M0.5), J160900.7-190852 (K7), [PZ99]J160357.6-203105 (K5), J155829.8-231007 (M3), and J161420.2-190648 (M0).
accretion disk models of Robitaille et al. (2006). To facilitate the observed SEDs are compared with those predicted by the substantial Upper Scorpius literature (e.g., Preibisch & Zinnecker 1999; Preibisch et al. 2002; Carpenter et al. 2006). A range of variability (Dahm & Carpenter 2009). To compare the observed SEDs with the accretion disk models, fluxes were measured in 0.8–70 μm SEDs are well sampled between 2.2 and 24 μm, a spectral region dominated by disk emission originating from the terrestrial region.

The Robitaille et al. (2006) grid of pre-computed, two-dimensional radiative transfer models consists of 20,000 young stellar objects in varying stages of evolution and viewed from 10 inclination angles. A total of 14 parameters are randomly sampled that specify stellar (e.g., M∗, R∗, T eff) as well as disk (e.g., M disk, M, R min) properties. The models have been successfully applied to the SEDs of 30 spatially resolved Class I and II sources in Taurus-Auriga by Robitaille et al. (2007), including the transition disk objects GM Aur and DM Tau.

The SEDs of nine Upper Scorpius members were compared to the accretion disk models. The sources excluded from the model fitting analysis were J160545.4-202308 and J155829.8-231007, the best-fitting, and maximum values of disk mass (M disk), and M, and inner disk radius (R min). The ranges of M disk and M are significant, in the model fitting analysis using the Paβ and Brγ emission line luminosities are consistent with those predicted by the best-fitting models.

As might be expected, inner disk radius appears to be better constrained by the subsets of best-fitting models than either M disk or M. Several Upper Scorpius sources, e.g., [PZ99]J161411.0-190648, [PZ99]J160357.6-203105, and ScoPMS 31, exhibit

| Source          | EW(Paβ) (Å) | EW(Brγ) (Å) | log L acc(Paβ)/L⊙ | log L acc(Brγ)/L⊙ | log M(Paβ)⊙ | log M(Brγ)⊙ |
|-----------------|------------|------------|--------------------|--------------------|-------------|-------------|
| [PZ99]J160357.6-203105 | −0.51      | −0.58      | −2.23 ± 0.18       | −2.00 ± 0.23       | −9.47 ± 0.18 | −9.23 ± 0.27 |
| J160000.7-190852  | −1.35      | −2.20      | −2.04 ± 0.15       | −1.74 ± 0.18       | −9.19 ± 0.15 | −8.89 ± 0.18  |
| J161420.2-190648  | −3.40      | −3.69      | −1.48 ± 0.10       | −0.65 ± 0.10       | −8.40 ± 0.10 | −7.60 ± 0.10  |
| ScoPMS 31         | ...        | −1.14      | −1.62 ± 0.17       | ...                | −8.52 ± 0.17 | ...          |
| J155829.8-231007  | −4.15      | −3.27      | −2.37 ± 0.20       | −2.50 ± 0.31       | −9.39 ± 0.20 | −9.52 ± 0.31  |

Notes.  
1 Negative equivalent width implies emission.  
2 Accretion luminosity determined using the linear relationship of Muzerolle et al. (1998).  
3 M derived assuming R∗, T eff, and M, values listed in Table 1.
ranges in $R_m$ from minimum to maximum of an order of magnitude or less. Five sources have best-fitting inner disk radii that exceed their respective dust sublimation radii: [PZ99]J161411.0-230536 ($R_m = 0.87$ AU), J160900.7-190852 (2.18 AU), ScoPMS 31 (9.62 AU), J161115.3-1917521 (0.29 AU), and J160357.9-194210 (2.0 AU for a typical $\sim 0.7 \, M_\odot$ pre-main-sequence star. The Upper Scorpius disks are presumably at an advanced evolutionary stage relative to those found around Class II sources in Taurus-Auriga and exhibit SEDs that are consistent with reduced levels of near and mid-infrared disk emission (Dahm & Carpenter 2009). Robitaille et al. (2007) find that all Taurus sources, except for the known transition disk objects (e.g., GM Aur, DM Tau) can be fit by models having disks and envelopes with inner disk radii equal to the dust sublimation limit. Approximately one-third of the Robitaille et al. (2006) models have inner disk radii set to the dust destruction radius. The remaining models have increasing inner disk radii that span from the dust destruction radius to 100 AU. The inner disk gaps are treated by the models as being completely evacuated of dust (Robitaille et al. 2006).

To constrain $R_m$ for the Upper Scorpius disk sample, a probability ($P$) is calculated for a given inner disk radius using the returned $\chi^2$ values from the model fits of the observed SEDs:

$$P = \frac{1}{\sqrt{2\pi}} e^{-\chi^2/2}.$$  \hspace{1cm} (4)

This probability density function assumes a normal distribution for the returned $\chi^2$ values of the individual fits. In the center panels of Figures 5(a)–(i), $P$ is plotted as a function of $R_m$ for each Upper Scorpius source. Superimposed in the figures as cross-hatched histograms are the distributions of inner disk radii for the entire sample of models considered for each source.

The current disk evolution scenario suggests that dust grains coagulate and settle toward the mid-plane prior to the formation of large planetesimals. The Robitaille et al. (2006) models use two disk structure parameters to mimic the effects of dust settling: a disk flaring parameter ($\beta$) and a disk scale height factor ($z$). If both of these parameters are low for a given model, this could be indicative of dust mid-plane settling (Robitaille et al. 2006). In the lower panels of Figure 5, $\beta$ is plotted as a function of $z$ for the entire sample of models considered for each source. Superimposed in red are the best-fitting models as determined by the $\chi^2 - \chi^2_{\text{min}} < N$ measure of goodness of fit. In general, the spread in both parameters for the best-fitting models is significant, suggesting that neither is well constrained. Some argument can be made that the SEDs of [PZ99]J161411.0-230536, J160643.8-190805, and ScoPMS 31 are better fit by models having lower than average values of $\beta$. The disk scale height factors for these sources, however, are found to vary significantly. Evidence for dust settling effects among the Upper Scorpius primordial disk sample remains inconclusive at best.

### Table 3

| Source               | $T_{\text{eff}}$ Range (K) | Number | Model ID | $A_V^d$ (mag) | $d_i^d$ (pc) | $\beta^d$ | $\chi^2_{\text{d}}$ |
|----------------------|-----------------------------|--------|----------|---------------|--------------|-----------|---------------------|
| [PZ99]J161411.0-230536 | 4900–5410                   | 500    | 3014849  | 1.0           | 125          | 75.5      | 6.30                |
| [PZ99]J160357.6-203105 | 4205–4590                   | 2993   | 3012326  | 2.0           | 125          | 31.8      | 13.44               |
| J160643.8-190805       | 4060–4350                   | 2753   | 3005191  | 0.5           | 185          | 75.5      | 1.51                |
| J160823.2-193001       | 4850–4060                   | 1190   | 3013235  | 0.75          | 165          | 41.4      | 0.34                |
| J160900.7-190852       | 4850–4060                   | 1190   | 3019185  | 1.0           | 145          | 81.4      | 4.72                |
| J161420.2-190648       | 3580–4060                   | 1998   | 3008376  | 3.0           | 125          | 56.6      | 58.91               |
| ScoPMS 31              | 3729–4060                   | 1917   | 3002397  | 2.0           | 165          | 81.4      | 2.19                |
| J161115.3-175721       | 3580–3850                   | 1409   | 3016046  | 2.0           | 165          | 41.4      | 2.87                |
| J160357.9-194210       | 3470–3720                   | 1180   | 3018769  | 0.5           | 155          | 63.3      | 0.42                |

**Notes.**

- $T_{\text{eff}}$ range considered for the adopted spectral type of the source.
- The number of models having $T_{\text{eff}}$ values within the specified range.
- The best-fitting model identification number from Robitaille et al. (2006).
- Extinction, distance, inclination angle, and reduced $\chi^2$ values for the best-fitting model.
Figure 5. Top: the observed 0.8–70.0 μm SEDs for nine late-type Upper Scorpius disk-bearing stars constructed from the SpeX 0.8–5.2 μm spectra, the Spitzer IRS 5.4–37.0 μm spectra, and the MIPS 24 and 70 μm photometry. Superimposed in red are the best-fitting model profiles of Robitaille et al. (2006), placed at their respective distances and extinctions. The disk contributions to the model SEDs are shown as dotted black curves, and the model stellar photospheres from the Kurucz (1979) atlas as solid blue curves. Center: the probability (P) of inner disk radii (R_{in}) for the Upper Scorpius disk-bearing sample plotted as a function R_{in}. Superimposed in the figures as cross-hatched histograms are the distributions of inner disk radii for the entire sample of models considered for each source. The vertical dashed lines represent the dust sublimation radii for the best-fitting models. Bottom: the disk flaring parameter (β) plotted as a function of disk scale height factor (z) for all models examined in the SED fitting process. The best-fitting models are shown in red.

(A color version of this figure is available in the online journal.)
Figure 5. (Continued)
Another limitation of the Robitaille et al. (2006) models is the assumption that disk gaps are completely devoid of dust, a reduction in complexity that impacts the predicted near- and mid-infrared excess distributions, critical for constraining disk emission originating from the inner disk rim.

In summary, three of nine late-type, disk-bearing stars in the Upper Scorpius sample exhibit SEDs that are most consistent with having inner disk radii that lie beyond the sublimation radius for silicate dust: [PZ99]J161411.0-230536, J160900.7-190852, and ScoPMS 31. The best-fitting models for two additional sources: J161115.3-175721 and J160357.9-194210, have inner disk radii that substantially exceed their respective dust sublimation radii, but the probability distributions of \( R_{\text{in}} \) values for these sources exhibit significant dispersion. Proposed indicators for the effects of mid-plane settling in the best-fitting Robitaille et al. (2007) models of the Upper Scorpius sample (i.e., concurrent decreased values of the disk flaring parameter and disk scale height factor) are inconclusive, leaving open to question the nature of the remaining Upper Scorpius primordial disks and their evolutionary state.

Transition disks are believed to represent an early stage of disk clearing and may be in the process of rapid grain growth, planetesimal and planet formation. Adopting the disk classification scheme of Luhman et al. (2010), many Upper Scorpius sources would be classified as pre-transitional objects, i.e., disks that exhibit reduced emission at wavelengths \( \lesssim 10 \mu m \), but that still retain significant disk emission at longer wavelengths. Two of the Upper Scorpius pre-transitional candidates appear to exhibit gapped disk structure: J160900.7-190852 and ScoPMS 31. Both sources are accreting (Dahm & Carpenter 2009 and Section 4) and both exhibit SEDs that are most consistent with having substantial inner disk radii, \( \sim 1–10 \) AU. At least one Upper Scorpius source, J160643.8-190805, exhibits an infrared SED that is suggestive of an homologously depleted disk system as defined by Currie et al. (2009).

Before attributing the reduced levels of infrared excess emission in the Upper Scorpius sample to disk evolutionary processes, binarity must be considered as an explanation for the cleared out inner cavities (e.g., CoKu Tau 4; Ireland & Kraus 2008). Two of the disk candidates included in this analysis, [PZ99]J161411.1-230536 and ScoPMS 31, are established binaries resolved by high angular resolution imaging. It must also be acknowledged that the SpeX sample represents only \( \sim 25\% \) of all primordial disk-bearing stars expected among the \( \sim 250 \) known late-type members of the Upper Scorpius OB association. Observations of more late-type, disk-bearing systems are critically needed to confirm the results presented here.

### Table 4

| Source               | \( M_{\text{Disk}} \) (M\(_{\odot}\)) | \( M \) (M\(_{\odot}\) yr\(^{-1}\)) | \( R_{\text{in}} \) (AU) |
|---------------------|-------------------------------------|-------------------------------------|---------------------|
|                     | (Min) | (Best) | (Max) | (Min) | (Best) | (Max) | (Min) | (Best) | (Max) |
| [PZ99]J161411.0-230536 | 7.18E-7 | 1.06E-5 | 1.16E-2 | 2.88E-13 | 2.43E-11 | 5.13E-8 | 0.10 | 0.87 | 2.24 |
| [PZ99]J160637.6-203105 | 2.60E-6 | 2.62E-5 | 3.47E-2 | 6.79E-13 | 1.16E-11 | 6.97E-8 | 0.06 | 0.07 | 0.14 |
| J160643.8-190805      | 5.30E-8 | 3.70E-6 | 5.34E-1 | 5.06E-15 | 2.28E-14 | 2.22E-9 | 0.04 | 0.05 | 10.00 |
| J160823.2-193001      | 6.07E-6 | 4.30E-4 | 1.72E-2 | 4.44E-12 | 2.02E-11 | 3.25E-8 | 0.03 | 0.04 | 1.19 |
| J160900.7-190852      | 1.44E-5 | 6.56E-3 | 6.56E-3 | 4.44E-12 | 1.43E-9 | 6.05E-9 | 0.04 | 2.18 | 4.13 |
| J161420.2-190648      | 4.28E-5 | 4.28E-5 | 1.36E-3 | 3.31E-10 | 3.31E-10 | 1.17E-7 | 0.08 | 0.80 | 0.80 |
| ScoPMS 31            | 5.90E-5 | 8.24E-4 | 1.45E-3 | 1.31E-12 | 3.57E-9 | 5.20E-9 | 4.08 | 9.62 | 15.70 |
| J161115.3-175721      | 3.52E-7 | 3.53E-5 | 1.78E-2 | 2.69E-13 | 1.74E-11 | 4.10E-8 | 0.03 | 0.29 | 1.19 |
| J160357.9-194210      | 3.62E-7 | 2.50E-4 | 9.50E-3 | 2.41E-13 | 2.16E-11 | 1.63E-8 | 0.02 | 0.17 | 1.29 |

\* Note. Minimum, best-fitting, and maximum \( M_{\text{disk}} \), \( M \), and \( R_{\text{in}} \) values of the subset of best-fitting models of Robitaille et al. (2006).
Improved modeling of the Upper Scorpius primordial disk-bearing sample is clearly needed to provide better constraints for the inner disk structure of these presumably evolved primordial disk systems. Monte Carlo three-dimensional, radiative transfer codes are now available that could be applied to the SEDs of these sources. High angular resolution imaging is also needed to identify close (≤10 AU) binary companions that may account for the reduced near- and mid-infrared excess emission of these sources relative to Class II sources in Taurus-Auriga. Infrared interferometric observations or precision radial velocity monitoring are also needed to identify tighter pairs or spectroscopic binaries that would be capable of dynamically clearing the disk interiors of these systems.

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