Three-Dimensional Metal-Insulator-Metal Decoupling Capacitors with Optimized ZrO₂ ALD Properties for Improved Electrical and Reliability Parameters

Konstantinos Efstathios Falidas 1,* , Kati Kühnel 1 , Matthias Rudolph 1 , Maximilian B. Everding 1, Malte Czernohorsky 1 and Johannes Heitmann 2

1 Fraunhofer Institute for Photonic Microsystems (IPMS), Center Nanoelectronic Technologies (CNT), An der Bartlake 5, 01109 Dresden, Germany
2 Institute of Applied Physics, TU Bergakademie Freiberg, Leipziger Str. 23, 09599 Freiberg, Germany
* Correspondence: konstantinos.efstathios.falidas@ipms.fraunhofer.de; Tel.: +49-351-2607-3071

Abstract: Embedded three-dimensional (3-D) metal-insulator-metal (MIM) decoupling capacitors with high-κ dielectric films of high capacitance and long-life time are increasingly needed on integrated chips. Towards achieving better electrical performance, there is a need for investigation into the influence of the variation in atomic layer deposition (ALD) parameters used for thin high-κ dielectric films (10 nm) made of Al₂O₃-doped ZrO₂. This variation should always be related to the structural uniformity, the electrical characteristics, and the electrical reliability of the capacitors. This paper discusses the influence of different Zr precursor pulse times per ALD cycle and deposition temperatures (283 °C/556 K and 303 °C/576 K) with respect to the capacitance density (C-V), voltage linearity and leakage current density (I-V). Moreover, the dielectric breakdown and TDDB characteristics are evaluated under a wide range of temperatures (223–423 K).

Keywords: high-κ; dielectric; ZrO₂; Al-doping; MIM capacitors; 3-D MIM capacitors; decoupling capacitors

1. Introduction

Placing a relatively large passive device between the supply rails of an integrated circuit is the most common technique to suppress power supply noise. Such devices, called decoupling capacitors, have been used both on- and off-chip for over 40 years. Their use enables improved noise immunity and consequently improve voltage regulation of the circuits. The metal-insulator-metal (MIM) capacitor is a capacitor widely used for decoupling purposes due to its high capacitance value while maintaining low parasitic capacitance and having low resistivity electrodes. In addition, MIM capacitors are able to deliver similar capacitance values to the multi-layer ceramic capacitors (MLCC) in much thinner layers. Therefore, MIM capacitors can be better integrated within the various levels of the chip fabrication process. As technology nodes advance to the nm regime, the existing MIM fabrication process needs to be further optimized in terms of capacitance density, reliability and space exploitation.

Capacitance densities may rise by reducing the dielectric thickness, which also achieves better device miniaturization for low voltage applications. However, this creates the problem of low breakdown voltage, poor voltage linearity and large leakage current. The introduction of dielectrics with higher dielectric constant than that of SiO₂ (>3.9), so-called high-κ dielectrics, overcome some of these issues. High-κ materials exhibit a larger permittivity and band gap; thus, they enhance the capacitance while limiting the leakage simultaneously. In planar form, decoupling capacitors with high-κ dielectrics deposited by atomic layer deposition (ALD) at low temperatures (<650 K) have been investigated for competing against thick silicon oxide MIM capacitors [1]. Devices with high-κ ZrO₂-based
dielectric have already been reported for decoupling purposes [2]. ZrO$_2$ has a wide band gap (5.1–7.8 eV) and a high dielectric constant, which varies from 17 to 47 depending strongly on its crystalline structure (monoclinic, cubic or tetragonal phase) [3,4].

To further increase the capacitance density of a MIM capacitor, while keeping the leakage current low, the shape of the capacitor has to change from a 2-dimensional to a 3-dimensional (3-D) structure. The increase in the surface-to-volume ratio for gaining higher capacitance densities has already been used in the past, for example in the case of DRAM cells with deep holes (trenches) etched in silicon or stacked oxide for miniaturizing the capacitors and increasing their storage [5]. Trenches with depth to width ratios (aspect ratios) of 50:1 are commonly used. It has been proven that the surface quality and the formation of the substrate in which the capacitor stack is deposited can strongly affect the final morphology of the capacitor [6], let alone in the case of 3-D structures where the density of the 3-D patterns as well as their geometry can further influence the electrical properties of the final product. Thus, conformal deposition and good uniformity of high-$\kappa$ dielectrics inside high aspect ratio 3-D structures is challenging even in the case of ALD processing.

The influence of important ALD parameters on the electrical properties of the dielectric has been analyzed in previous studies in terms of the growth rate of the metal precursor or the variation in the oxidation pulse during the deposition of ZrO$_2$-based 3-D MIM capacitors [7,8]. However, there are further ALD parameters, which affect the performance of the ALD process including the precursor pulse times and the deposition temperature which need to be understood to define the appropriate recipes for achieving a conformal deposition with low carbon content and proper sidewall coverage within the deep trenches of the 3-D MIM decoupling capacitors. Additionally, it has been shown that introducing Al with its large bandgap of 8.7 eV and its low dielectric constant of ~9 within the ZrO$_2$-based high-$\kappa$ stack keeps the dielectric amorphous during deposition, avoiding unwanted premature crystallization [2,8]. Consequently, high electric field strength, low oxygen diffusivity, lifetime improvement and lower leakage currents can be achieved [9,10]. Aluminum can be deposited with the same deposition technique (ALD) and it can be combined with the deposition of the high-$\kappa$ dielectric forming laminated or nano-laminated/mixed stacks. Specific capacitances of ~27 nF/mm$^2$ have been reported for planar capacitors including nano-laminates of Al$_2$O$_3$ in ZrO$_2$ [11].

The electrical data of the decoupling capacitors have to still fulfill the industrial specifications which dictate, for example, that the leakage current has to remain as low as possible (<1 µA/cm$^2$) and the stationary capacitance over applied voltage (capacitance voltage linearity) has to remain consistent during the whole operating lifetime of the product to warrant the reliability of the capacitor. Moreover, long-term lifetime under various operating temperatures, regardless the size of the decoupling capacitor plays an essential role in all stages of product development. Temperature- and time-dependent dielectric breakdown (T-TDDB) should also meet industrial reliability requirements, which note that a minimum ten-year lifetime criterion at accumulated fail rate of 10 ppm should be satisfied. In this context, the appropriate process parameters, especially for the ALD process (Zr precursor pulse times, deposition temperature) of the Al$_2$O$_3$-doped ZrO$_2$ high-dielectric stack have to be established, enabling optimum electrical properties.

In this work, we successfully fabricated 3-D MIM capacitors with an ALD deposited Al$_2$O$_3$-doped ZrO$_2$ conformal and smooth high-$\kappa$ dielectric stack with limited impurities on high aspect ratio substrates. We investigated the influence of the high-$\kappa$ metal (Zr) precursor pulse time and the deposition temperature of the dielectric layers on the uniform coverage of the 3-D capacitors. The uniformity is a key performance factor as it is inextricably linked to the electric performance of the devices. We evaluated the fabricated devices experimentally with respect to electrical properties (capacitance and current dependence on voltage) and reliability parameters (TDBB).
2. Materials and Methods

The 3-D MIM capacitors (Figure 1a,d) were formed on highly boron-doped (0.02 Ω·cm), 775 µm thick, 300 mm silicon (100) wafers. In order to create the deep trenches needed for the 3-D shape of the capacitors, the wafers were structured using i-Line photolithography on a SiO₂ hard-mask. Test devices with different footprint areas (0.01–0.1 mm²) and trench structures with depth of ~7.7 µm, as shown in Figure 1a, were defined and could be realized.

Table 1. Experimental conditions and thickness values of Al₂O₃-doped ZrO₂ (ZrAlₓOᵧ) with average thickness of ~23 nm (see in Figure 1b,c) and Figure 2b were deposited by atomic layer deposition (Jusung Eureka 3000, warm wall reaction chamber) using TEMAZr (Tetrakis-ethylmethylamino-zirconium, Zr[N(C₂H₅)₂CH₃]₄) and TMA (Tri-methyl-aluminum, Al(CH₃)₃) as metal precursors (353 K precursor bubbler temperature), ozone (O₃) with flow rate of 200 g/m³ as reactant/oxidizing agent and Argon (Ar) as carrier/purge gas at deposition temperatures of 556 K or 576 K. The latter temperature has been reported to be the maximum allowed ALD temperature, before decomposition of TEMAZr begins [14]. The TEMAZr cycle consisted of pulsing with TEMAZr, purging with Ar gas, oxidizing with ozone, and purging with Ar gas. In a similar way, the dopant (TMA) was injected, followed by purging steps with Ar, O₃ pulse and Ar purge again. In Table 1, the physical thickness values are presented which were measured on the top lateral surface of the high-k film stack. TEMAZr pulse time was increased every time while decreasing the number of deposition cycles per super-cycle. Oxidation time (O₃ pulse duration) followed the changes of the precursor pulse time.

![Figure 1. (a) SEM cross-section of a 7.7 µm deep array of trenches; (b) detail of one of the trenches in (a) at the top trench section; (c) detail of one of the trenches in (a) at the bottom of the trench; (d) schematic cross-section (not to scale) of the 3-D MIM capacitor realized in a silicon substrate. Bottom contact can be contacted on the front side or on the rear side.](image-url)
were changed following the variation in the TEMAZr pulse times in order to not affect the ALD temperature of 576 K and Zr precursor pulse times of 12 s/cycle:

Table 1. Experimental conditions and thickness values of Al electrode layers and the dielectric stack (ZrAlxOy) are indicated.

| TEMAZr Pulse Time (t_pulse) | TEMAZr Cycles/Super-Cycle | Deposition Temperature (T_D) | ZrAlxOy Physical Thickness ¹ |
|----------------------------|---------------------------|-----------------------------|-----------------------------|
| 3 s                        | 61                        | 283 °C/556 K                | -                           |
| 6 s                        | 59                        | 283 °C/556 K                | 24.6 nm                     |
| 9 s                        | 59                        | 283 °C/556 K                | 28.7 nm                     |
| 12 s                       | 55                        | 283 °C/556 K                | 29.1 nm                     |
| 12 s                       | 48                        | 303 °C/576 K                | 27.4 nm                     |
| 15 s                       | 48                        | 303 °C/576 K                | 27.0 nm                     |

¹ measured on the top lateral surface of the high-κ film stack.

After 48–61 cycles of TEMAZr and O₃, ALD cycles of TMA with O₃ followed, so that a thin film of Al₂O₃ was formed on top of the ZrO₂ stack. The aluminum ALD cycles were changed following the variation in the TEMAZr pulse times in order to not affect the stoichiometry between Al and Zr in the dielectric stack. A series of ZrO₂ cycles (48–61, depending on the pulse time of TEMAZr) followed by Al₂O₃ cycles defined a so-called super-cycle. By repeating this supercycle of ZrO₂ and Al₂O₃ four times, the desired film thickness of the high-κ dielectric stack was achieved, with aluminum distributed across the dielectric stack. The last super-cycle consisted only of the TEMAZr cycle in order to have ZrO₂ as the top layer of the high-κ dielectric. There was no post deposition anneal (PDA) in order to keep the high-κ films in the amorphous state with a low defect density.

Figure 2. TEM micrographs of one individual trench of the deposited MIM stack with increased ALD temperature of 576 K and Zr precursor pulse times of 12 s/cycle: (a) bottom area of the 7.7 µm deep trench; (b) details of the sidewall of a trench with the MIM stack visible, where individual TiN electrode layers and the dielectric stack (ZrAlₓOᵧ) are indicated.

A second photolithographic patterning step with a boron-doped silicon electrode was used to define the top electrode of the devices and served also as filling of the void of the trenches, while keeping the physical series resistance of the electrodes and therefore the Equivalent Series Resistance (ESR) as low as possible.

The dependence of the capacitance density on temperature for increasing TEMAZr pulse times can be checked in extreme operating environmental conditions, similarly to those can be checked in extreme operating environmental conditions, similarly to those...
The morphology of the MIM capacitors were characterized by Scanning Electron Microscopy (S-5000, Hitachi). Samples were cleaved just before being loaded into the SEM chamber to maintain a fresh cross-sectional surface. Additionally, the MIM stack was visually analyzed using Transmission Electron Microscopy (F20 TEM 200 kV, FEI Tecnai) to support the SEM findings and additionally investigate more the high-κ dielectric formation in the trench structures. TEM was also used to investigate the crystal status of the fabricated samples after the end of the integration process.

The electrical characterization was performed on a temperature-controlled wafer stage of a fully automated semiconductor probe station (Precio Nano, Tokyo Electron) at room temperature (298 K). Capacitance (C-V) and leakage current (J-V) characteristics were determined using a C-V analyzer (E4980A, Agilent/Keysight) and a semiconductor parameter analyzer (B1500, Agilent/Keysight), respectively. In this work, a series equivalent circuit model was used for the capacitance measurements and voltage sweeps from negative to positive polarity directions were measured. As a way to have a high statistic over the wafer, all electrical measurements were executed on capacitors with a top lateral surface area of 0.01 mm² and at minimum 53 dies per wafer. Throughout this work and more specifically in Figures 3a, 5a–c, 7a and 9, box-charts are used for the graphical representation of the key statistical values of the previously mentioned datasets of min. 53 capacitors. In particular, each separate set of data is represented as a separate box. The box is determined by the 25th and 75th percentiles. The whiskers are determined by the 5th and the 95th percentiles. The horizontal line in each box depicts the median and the open square symbol the mean, respectively. The samples were kept in the dark during electrical characterization. The area enhancement factor of the capacitors with high aspect-ratio holes compared to planar capacitors with similar surface area has been found to be ~20 in a previous experiment with the same setup and layout [17]. Within this work, the samples have always been normalized over the top lateral surface area of each capacitor.

3. Results and Discussion

3.1. Capacitance Behavior, Temperature Dependence and Voltage Linearity

Figure 3a depicts the results of the capacitance-voltage measurements at zero bias voltage \(C_0\) and 10 kHz of the experiments of Table 1. The capacitance density values slightly increased with increasing TEMAZr pulse time. At 3 s pulse time and \(T_D = 556\) K, the \(C_0\) reached \(\sim 260 \text{nF/mm}^2\) but with a very high deviation (±25%) over the wafer (non-uniformity). Due to the high non-uniformity and the high number of impurities, a parasitic parallel resistance appeared, increasing the equivalent series resistance (ESR) of the sample remarkably (~750 Ω), thus indicating a very high leakage (Figure 3b). By increasing the TEMAZr pulse time to 6 s and higher, the ALD process reached its saturation point, which can be verified by the significant lower deviation of the \(C_0\) values and the drop of the ESR values (Figure 3). Despite reaching the ALD saturation point, higher TEMAZr pulse times (9–12 s) were tested while aiming for a homogenous sidewall coverage all over the 3-D trench structure (bottom-middle-top). In such case, there was only a minimal increase in \(C_0\), which came together with slightly higher deviation of the values.

At the same time, the increase in the deposition temperature by 20 K (\(T_D = 576\) K) while keeping the TEMAZr pulse time constant (12 s) improved the uniformity of \(C_0\) (see also Figure 4a) and slightly increased its value. A further increase of the pulse time to 15 s showed similar behavior to the elevated pulse times of 9 s and 12 s for \(T_D = 556\) K, namely only a minimal rise in \(C_0\) with higher non-uniformity. The small thickness variations both inside the trench holes (Figure 2) as well as on the top surface of the capacitors have not significantly affected the capacitance behavior, but only improved the sidewall coverage.
The dependence of the capacitance density on temperature for increasing TEMAZr pulse time and deposition temperature was studied under a wide temperature range (223–423 K), as depicted in Figure 5a–c. From these outcomes, the performance of capacitance can be checked in extreme operating environmental conditions, similarly to those expected for automotive applications [18]. In all cases, there is a temperature acceleration and the capacitance density at zero bias ($C_0$) was found to vary in a fairly linear manner with respect to temperature. By increasing the TEMAZr pulse time from 6 s to 12 s, negligible changes appear in the $C_0$/T slope (~4.2 nF/mm² K). On the other hand, by increasing the deposition temperature from 556 K to 576 K, the temperature acceleration rate doubles (~8.7 nF/mm² K). The majority of the defects seem to be caused during the ozone steps at higher deposition temperature. These defects are activated by temperature and start acting as traps. By that, more charges can reach the electrodes faster while the rate of applied voltage remains constant. Thus, the capacitance density increases based on the standard equation of $dq/dt = CdV/dt$. 

**Figure 3.** Variation over different TEMAZr pulse times and high-κ deposition temperatures measured at 0 V, 10 kHz and room temperature (298 K) for (a) capacitance density ($C_0$) and (b) the median of equivalent series resistance (ESR).

**Figure 4.** Wafer−maps of data measured at room temperature (298 K) on all (523) wafer−dies of a ZrAl$_x$O$_y$ dielectric stack with 12 s TEMAZr pulse time and 573 K deposition temperature: (a) $C_0$ measured at frequency of 10 kHz; (b) leakage current density at +3 V plotted in logarithmic scale.
Figure 5. $C_0$ over temperature (223 K up to 423 K), at measurement frequency of 10 kHz for: (a) 6 s TEMAZr pulse time and $T_D = 556$ K, (b) 12 s TEMAZr pulse time and $T_D = 556$ K deposition temperature and (c) 12 s TEMAZr pulse time and $T_D = 576$ K; (d) $C(V)$—plot over temperature of the ZrAl$_x$O$_y$ stack of (c).

The temperature coefficient of capacitance $\alpha_T$, a figure of merit of the dielectric capacitor devices, which describes the maximum change in capacitance density over a specified temperature range, is formulated as

$$\alpha_T = \frac{C_T}{C_{RT}} - \frac{1}{T_{RT}} \cdot 10^6 \text{ppm/K}$$

(1)

where $C_T$ and $C_{RT}$ are the capacitance values at a particular temperature limit T and at room temperature ($T_{RT} = 298$ K), respectively. Room temperature is established as a reference temperature. The numerical values for 12 s TEMAZr pulse time and $T_D = 576$ K were determined to be 322 ppm/K and 280 ppm/K at the extreme temperatures of 223 and 423 K, respectively and are within the limit of 750 ppm/K set by the Electronic Industries Alliance (EIA) for capacitors [19].
An important parameter of a MIM capacitor for decoupling applications is its capacitance stability over bias voltage—also called voltage linearity. The voltage linearity of MIM capacitors is determined by a second-order polynomial equation

\[
\frac{C(V)}{C_0} - 1 = \alpha V^2 + \beta V
\]

where \(V\) is the applied bias voltage; \(C(V)\), the capacitance density at a particular bias voltage; \(C_0\) the zero-bias capacitance; \(\alpha\) and \(\beta\), the quadratic and linear Voltage Capacitor Coefficients (VCC), expressed in units of parts per million (ppm)/\(V^2\) and ppm/\(V\), respectively.

The capacitance density values for the ZrAl\(_x\)O\(_y\) stack with TEMAZr pulse time of 12 s and \(T_D = 576\) K were measured at a measurement frequency of 10 kHz from \(-3\) V to \(+3\) V with a voltage step of \(+0.2\) V at temperatures from 223 K to 423 K and are depicted in Figure 5d. The data was fitted with the parabolic curve of Equation (2). The extracted quadratic term \(\alpha\) exhibits a linear relationship with respect to the applied temperature (see the inset of Figure 5d) which can be expected from the electrode polarization model and explained by thermally enhanced mobility of charge carriers in the dielectric [6].

Both the temperature coefficient of capacitance \(\alpha_T\) (1) and the voltage capacitance coefficients, \(\alpha\) and \(\beta\) (2) are independent of the area enhancement factor of the 3-D capacitors.

Another important parameter, which defines the quality of the decoupling capacitors, is their performance over different applied frequencies. Figure 6 shows the plot of capacitance density values measured from \(-3\) V to \(+3\) V at different frequencies (10–100 kHz) for 12 s TEMAZr pulse time and \(T_D = 576\) K. On the one hand, the capacitance density remained nearly stable over voltage at any fixed applied frequency and voltage polarity, which indicates the good capacitor’s stability under a continuously increasing voltage stress. On the other hand, by increasing the measurement frequency, the capacitance was notably reduced. Because of the increasing measurement frequencies, the polarization density \(P(t) = P_0 + \int_0^t P(t)dt\) cannot fully respond due to dielectric relaxation. Consequently, the capacitance density drops, regardless the amount of the applied bias voltage. In the case of 3-D decoupling capacitors with high aspect ratio structures, this behavior becomes more noticeable due to their complex morphologies, such as edges, corners, and sidewall spikes, resulting from the etching process. These cause the localization of enhanced electric field which might further delay the response of charge and thus cause a reduction in the quadratic VCC, as depicted in Figure 6b, where the extracted \(\alpha\) values decreased from 2300 ppm/\(V^2\) to 10 ppm/\(V^2\) with increasing frequency.

### 3.2. Leakage Currents, Breakdown and Reliability Characteristics

The leakage current density of the 3-D MIM capacitors is a critical factor in terms of reliability, especially for the very thin embedded MIM capacitors. The plot in Figure 7a reveals very low values (<1 \(\mu\)A/cm\(^2\)) for all tested material stacks at room temperature (298 K) and \(+3\) V bias voltage. With increasing the deposition temperature (\(T_D = 576\) K) of the high-k dielectric stack, the uniformity of the leakage current data increased. The increase in the deposition temperature together with a higher TEMAZr pulse time proved to be the most beneficial as it allowed capacitors not only to have an optimal low leakage current, but also the overall highest capacitance, which can be extracted from the I/C-plot in Figure 7b. Additionally to the improved uniformity of the leakage current over the wafer (Figure 4b), samples obtained with equal TEMAZr pulse time (12 s) but higher \(T_D\), demonstrated absolute values of leakage current density, which were decreased by almost two orders of magnitude. This large reduction can be related to both lower carbon impurities during the ALD deposition process and the lightly better step coverage at higher temperatures.
3.2. Leakage Currents, Breakdown and Reliability Characteristics

The leakage current raised with increasing measurement temperatures. This increase can be related to the thermally enhanced mobility of charge carriers. It has been reported before [6] that in case of amorphous Al-doped ZrO₂ high-κ thin films with a dielectric thickness higher than 10 nm, Poole-Frenkel electron emission from traps in the field region, which is in line with what was previously published [6].

Additionally, the leakage current behavior of the sample with the lowest leakage has been characterized under temperature for both voltage polarities (Figure 8a). An almost perfect symmetry of the leakage current density curves for both polarities could be extracted. The leakage current raised with increasing measurement temperatures. This increase can be related to the thermally enhanced mobility of charge carriers. It has been reported before [6] that in case of amorphous Al-doped ZrO₂ high-κ thin films with a...
dielectric thickness higher than 10 nm, Poole-Frenkel electron emission from traps in the dielectric is the dominant conduction mechanism due to the significant temperature dependence of the J-E characteristics. Additionally, the symmetry of the leakage current density curves for both polarities is a strong indicator both for non-fully crystalline dielectric and for Poole-Frenkel emission in case of MIM capacitors using metals with different work functions [20,21]. The Poole–Frenkel effect describes the way electrons in insulators use the thermal fluctuations to be promoted to the conduction band to travel between traps. The applied field lowers the barrier of the trap; hence a smaller work function is needed to overcome it [22]. Having that in mind and to further understand the conduction mechanism of 3-D MIM decoupling capacitors at low and high electric field strengths in both polarities, the J-E plot was re-plotted in Poole-Frenkel coordinates (Figure 8b) and the values were fitted with the Poole-Frenkel (PF) acceptor limiting conduction current:

$$J_{PF} \propto E \frac{e^{-\frac{\varphi_{PF}}{kT}}}{E^{\frac{\sqrt{\varepsilon_{opt}qE}}{\pi\varepsilon_0}} \frac{\sqrt{\varepsilon_{opt}qE}}{\pi\varepsilon_0}}$$

where $\varphi_{PF}$ describes the trap depth without external E-field, $\varepsilon_{opt}$ denotes the optical dielectric constant, q is the electron-charge, E stands for the electrical field applied to the film, $k_B$ is the Boltzmann’s constant and T is the applied temperature [21]. The PF-plot can be well-fitted in both polarities by a straight line up to very low field values of high electric field region, which is in line with what was previously published [6].

![Figure 8. (a) Leakage current density over the electric field at measuring temperatures from 223 K to 423 K; (b) Poole–Frenkel coordinates over temperature for both field polarities. Both data refer to the samples with 12 s TEMA Zr pulse time and $T_D = 556$ K.](image)

After measuring the leakage current density within the operating voltage range of the MIM decoupling capacitors (±3 V), breakdown voltage values and Time Dependent Dielectric Breakdown (TDDB) were also measured at temperatures from 223 K to 423 K on wafer level, using a voltage ramp stress (VRS) and constant voltage stress (CVS), respectively. To enable the comparison of different stacks/ALD parameters with different thickness values, the breakdown electric field strength was used, calculated according to the standard equation of $E = V/d$, assuming there is no electron shielding due to the absence of ferroelectricity [23]. Measurement of breakdown for 523 dies (full map) yielded values between 2.05 MV/cm and 2.65 MV/cm (Figure 9) with no early breakdown registered.
This indicates a low defect density because of the good quality of the deposition on the deep-trench structures. By increasing TEMAZr pulse times, similarly to the behavior of the leakage current density, more impurities caused lower breakdown fields despite the better trench-sidewall coverage. As an example, 6 s higher pulse time caused a drop of the breakdown field by 30%. Higher TD increased the breakdown of the films again to higher values.

![Figure 9. Breakdown Field at room temperature for increasing TEMAZr pulse time and deposition temperatures.](image)

The trade-off between the capacitance, the leakage current and the breakdown field, the standard electrical characterization parameters of a MIM capacitor, as depicted in Figures 7b and 9 indicated that the best ALD recipe proved to be the one with 12 s TEMAZr pulse time and higher TD.

The high breakdown field strength generally enables capacitors to have a large and stable working voltage and reflects the capacitor’s lifetime. In order to analyze the reliability of the 3-D MIM decoupling capacitors, the stress time required for dielectric breakdown under constant field stress was measured which is also called TDDB lifetime. For the TDDB measurements, the capacitors are subject to field stress in a range of 1.8–2.6 MV/cm. For each field stress value, multiple capacitors, each on an individual die, are assessed according to previously measured breakdown voltage and breakdown current density characteristics. The time to breakdown (t_{BD}) is measured individually with a sampling interval of 100 ms and a current compliance of 1 mA. From the measured t_{BD} data the breakdown time at 63.2% failure of the devices (t_{63.2}) is extracted, which for the purpose of this paper will be called characteristic lifetime. The characteristic lifetime data is then extrapolated to the operating electric field by fitting the data to the “linear E” model and calculating the voltage acceleration γ. This model was selected because it describes the metal-oxide bond breakage [24] while also being a more conservative model compared to the power law model used before for similar devices [25].

As shown in Figure 10, the extrapolation of the characteristic lifetime to the operating electric field (voltage) of ~1.2 MV/cm (3 V) yielded average values of >10^{10} s or more than 1000 years for samples with 12 s Zr precursor pulse time and an ALD deposition temperature of 576 K. This performance was confirmed for all measuring temperatures (T_m) up to 323 K. The characteristic lifetime of the 3-D MIM decoupling capacitors under room temperature conditions (298 K) has been improved by almost two orders of magnitude by increasing TD by 20 K due to the reduction of impurities.
Figure 9. Breakdown Field at room temperature for increasing TEMAZr pulse time and deposition temperatures (553 K, 573 K) for different stress fields measured at 273 K, 298 K and 323 K. Lifetime extrapolated to 10 years of lifetime using an exponential relation with voltage acceleration γ being the slope.

4. Conclusions

Embedded 3-D MIM decoupling capacitors with Al₂O₃-doped ZrO₂ high-κ dielectric film stack have been realized on silicon wafers. The behavior of the samples with respect to the complex morphologies of the 3-D trenches and ALD deposition conditions has been analyzed. Their electrical and reliability properties were investigated by capacitance, leakage current, dielectric breakdown and TDDB measurements. We have shown that with increased ALD deposition temperature (576 K) and longer Zr precursor pulse times (12 s/cycle) close to the precursor saturation point, the capacitors revealed a very high uniformity (>90%). Additionally, capacitance density with average values of 350 nF/mm² and low leakage current of <1 µA/cm² at +1.1 MV/cm (~+3 V) were measured. In addition, the leakage current behavior over temperature has been theoretically examined and can be attributed to the Poole-Frenkel effect. The breakdown field reached a maximum of ±2.7 MV/cm (±8 V). Finally, a characteristic lifetime of more than 10 years at the operating voltage of 3 V could be achieved successfully. These results, which meet the performance standards required from the semiconductor industry, highlight the importance of the careful tuning of the metal precursor and the ALD deposition temperature in high-κ dielectric stack systems used for 3-D decoupling MIM capacitors. Moreover, they emphasize the need for the further understanding of the interactions of ALD processes for such thin ZrAlₓOᵧ films. With that understanding, the material system can be further improved and become a more reliable candidate for embedded decoupling 3-D MIM capacitors.

Author Contributions: Conceptualization, K.E.F.; methodology, K.E.F.; software, K.E.F.; validation, K.E.F.; formal analysis, K.E.F. and K.K.; investigation, K.E.F.; resources, K.E.F., K.K. and M.R.; data curation, K.E.F.; writing—original draft preparation, K.E.F.; writing—review and editing, K.E.F., M.B.E., M.C. and J.H.; visualization, K.E.F.; supervision, M.C. and J.H.; project administration, M.C.; funding acquisition, M.C. All authors have read and agreed to the published version of the manuscript.
Funding: This research was funded via subcontract from Globalfoundries Dresden Module One within the framework Important Project of Common European Interest (IPCEI), by the German Federal Ministry for Economics and Energy (BMWi) and by the State of Saxony.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: Research data in this article is not shared.

Acknowledgments: The authors would like to thank Katrin Zimmermann and Kati Biedermann, who provided us with the SEM and TEM micrographs.

Conflicts of Interest: The authors declare no conflict of interest.

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