A NEW METHOD FOR THE ASSESSMENT OF AGE AND AGE SPREAD OF PRE-MAIN-SEQUENCE STARS IN YOUNG STELLAR ASSOCIATIONS OF THE MAGELLANIC CLOUDS\textsuperscript{*}.\textsuperscript{†}

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ABSTRACT

We present a new method for the evaluation of the age and age spread among pre-main-sequence (PMS) stars in star-forming regions in the Magellanic Clouds, accounting simultaneously for photometric errors, unresolved binarity, differential extinction, stellar variability, accretion, and crowding. The application of the method is performed with the statistical construction of synthetic color–magnitude diagrams (CMDs) using isochrones from two families of PMS evolutionary models. We convert each isochrone into two-dimensional probability distributions of artificial PMS stars in the CMD by applying the aforementioned biases that dislocate these stars from their original CMD positions. A maximum-likelihood technique is then applied to derive the probability for each observed star to have a certain age as well as the best age for the entire cluster. We apply our method to the photometric catalog of \(~\)2000 PMS stars in the young association LH 95 in the Large Magellanic Cloud, based on the deepest HST/ACS imaging ever performed toward this galaxy, with a detection limit of \(V \sim 28\), corresponding to \(M \sim 0.2 M_\odot\). We assume the initial mass function and reddening distribution for the system, as they have been previously derived by us. Our treatment shows that the age determination is very sensitive to the considered grid of evolutionary models and the assumed binary fraction. The age of LH 95 is found to vary from 2.8 Myr to 4.4 Myr, depending on these factors. We evaluate the accuracy of our age estimation and find that the method is fairly accurate in the PMS regime, while the precision of the measurement of the age is lower at higher luminosities. Our analysis allows us to disentangle a real age spread from the apparent CMD broadening caused by the physical and observational biases. We find that LH 95 hosts an age spread that is represented well by a Gaussian distribution with an FWHM of the order of 2.8–4.4 Myr depending on the model and binary fraction. We detect a dependence of the average age of the system with the stellar mass. This dependence does not appear to have any physical meaning, being rather due to imperfections of the PMS evolutionary models, which tend to predict lower ages for the intermediate-mass and higher ages for the low-mass stars.

Key words: Hertzsprung–Russell and C–M diagrams – methods: statistical – open clusters and associations: individual (LH 95) – stars: formation – stars: pre-main sequence

Online-only material: color figures

1. INTRODUCTION

The Large and Small Magellanic Clouds (LMC and SMC), the dwarf metal-poor companion galaxies to our own, are well-suited environments for the study of resolved star formation at low metallicity. Indeed, both galaxies are characterized by an exceptional sample of star-forming regions, where hydrogen has been already started to be ionized by the UV winds of massive stars, the H\textsc{ii} regions (Henize 1956; Davies et al. 1976). A plethora of young stellar associations and clusters are embedded in these regions (e.g., Bica & Schmitt 1995; Bica et al. 1999), giving evidence of recent (or current) clustered star formation. For many years, studies of the bright stellar content of these systems have provided and still provide unprecedented information on massive stellar evolution at low metallicities (see, e.g., Walborn et al. 1995, 2010). On the other hand, their faint stellar content, which is still in the pre-main-sequence (PMS) phase, became only recently accessible, exclusively through imaging with the Hubble Space Telescope (HST; see e.g., Gouliermis et al. 2006; Nota et al. 2006).

The identification of the low-mass PMS stars of star-forming regions in the Milky Way is quite demanding, requiring various techniques, such as variability surveys, spectroscopy, the detection of emission lines, and excess due to accretion (e.g., Briceño et al. 2007), in order to decontaminate the observed stellar samples from the evolved populations of the Galactic disk. This is not a problem for the observational study of the low-mass population of young clusters in the Magellanic Clouds (MCs), where memberships can be assigned simply from the position of the stars in the color–magnitude diagram (CMD) alone. This is guaranteed by the negligible distance spread of the LMC or SMC field, in comparison to the distance of the two galaxies.

Naturally, this observational advantage has important implications in our general understanding of clustered star formation and in consequence of PMS evolution, the initial mass function (IMF) and dynamical evolution of young star-forming clusters. Clearly, in order to detect low-mass stars in the MCs, imaging with high resolution and high sensitivity in optical wavelengths is required, and the Advanced Camera for Surveys (ACS) on board the HST has been so far an ideal instrument for this purpose. Indeed, deep imaging at a detection limit of \(m_{555} \simeq 28\) with ACS/WFC of the young LMC cluster LH 95 (Lucke &
Hodge 1970) proved for the first time that young MC stellar systems host exceptionally large numbers of faint PMS stars (Gouliermis et al. 2007), providing the necessary stellar numbers to address the issue of the low-mass IMF in star-forming regions of another galaxy (Da Rio et al. 2009).

PMS stars remain in the contraction phase for several times 10 Myr, and therefore they are chronometers of star formation and its duration. However, the estimate of stellar ages in the PMS regime from photometry is not without difficulties. This is because the observed fluxes are affected by several physical and observational effects. While these stars are optically visible, they are still surrounded by gas and dust from their formation, typically in circumstellar disks. Irregular accretion, due to instabilities in these disks, produces variability in the brightness of these stars. Moreover, PMS sources exhibit periodic fluctuations in light produced by rotating star spots. The variable nature of PMS stars dislocates them from their theoretical positions in the CMD. This effect in addition to observable biases, such as crowding and unresolved binarity, results in a broadening of the sequence of PMS stars in the CMD. As a consequence, the measurement of the age of the host system and any intrinsic age spread, which may indicate continuous star formation, cannot be derived from simple comparisons of PMS isochrones with the observed CMD. This problem is particularly critical for our understanding of the star formation processes in every region. In particular, Naylor (2009) showed that neglecting unresolved binaries may lead to an underestimation of the age of a young cluster by a factor of 1.5–2. Concerning the star formation duration, the apparent luminosity spread in the CMDs, if not disentangled from a possible intrinsic age spread, leads to an overestimation of the star formation timescale.

With spectroscopic and variability surveys with HST in star-forming regions of the MCs being practically impossible, the only way to disentangle the age and the duration of star formation in such regions from the rich samples of PMS stars is with the application of a statistical analysis, from a comparison between the observed CMDs and simulated ones.

In this study, we present a new maximum-likelihood method, for the derivation of the ages of PMS stars and the identification of any age spread among them. We apply our method to the rich sample of PMS stars detected with ACS in LH 95 to determine the age of the stars from their positions in the CMD in a probabilistic manner, accounting for the physical and observational biases that affect their colors and magnitudes. Our data set is presented in Section 2, and the evolutionary models to be used for our synthetic photometry are described in Section 3. In order to achieve the most accurate results, we further refine our previous selection of PMS star members of the system in Section 4. In Section 5, we construct modeled two-dimensional density distributions of PMS stars in synthetic CMDs by applying the considered observational and physical biases to the evolutionary models, and we apply our maximum-likelihood technique for the derivation of the age of the system. In Section 6, we use our method to investigate the existence of an age spread in LH 95, and we discuss our findings in Section 7. The summary of this study is given in Section 8.

2. THE DATA SET

In this study, we make use of the photometry of the young association LH 95, obtained from deep HST imaging with the Wide-Field Channel (WFC) of the ACS in the filters F555W and F814W, roughly corresponding to standard V and I bands, respectively (program GO-10566; PI: D. Gouliermis). The reduction of the images and their photometry is discussed in Gouliermis et al. (2007) and is presented in its entirety by Da Rio et al. (2009, hereafter Paper I). The photometric completeness reaches the limit of 50% at $m_{555} \approx 27.75$ mag and $m_{814} \approx 27$ mag, reduced due to crowding by $\sim 0.5$ mag in the central area occupied by the association. A control field $\sim 2'$ away from the system was also observed with an identical configuration. It has been used in Paper I to isolate the true stellar population of the association from that of the general LMC field by performing a statistical field subtraction on the observed CMD. The HST observations of the two fields and in the two bands were not consecutive. In particular, the I-band imaging of LH 95 was carried out four days after the V band. This will be crucial in the analysis we present in this work to model the observed effect of rotational variability.

The derived cluster-member population includes 2104 stars located in the PMS and upper main-sequence (UMS) region of the CMD with corresponding masses as low as $\sim 0.2$ $M_\odot$. In Paper I, we isolate the central area of the observed field of view (FOV), which covers a surface of 1.43 arcmin$^2$ ($\sim 260$ pc$^2$) corresponding to about 12.6% of the whole FOV, as the most representative of the association LH 95. This area, which encompass the vast majority of the observed PMS population, is referred to as the “central region.”

Given the low density of stars in LH 95 (of the order of 10 stars pc$^{-2}$ down to 0.3–0.2 $M_\odot$), it can be considered a loose association. In our previous work, three denser substructures—referred to as “subcluster A, B, C”—were identified. They include about half of the stellar population, while another half is distributed within the central region. In Paper I, we roughly estimated the dynamical properties of the central region as well as the three subclusters. This was done by estimating the total stellar mass, extrapolating our IMF for masses below our detection limit, and assuming a velocity distribution for the stars of the order of $\sim 2$ km s$^{-1}$. For the entire region, the crossing time is comparable with the age of the system, and it is few times shorter when computed singularly for each of the three subclusters. It was still unclear whether these stellar agglomerates were formed independently on each other, from separate star-forming events, or if their spatial distribution was simply caused by a clumpy distribution of the ISM in the parental cloud, during a single star formation event. In Section 5.3, we investigate this, excluding the presence of a spatial variability of the stellar ages.

3. DESCRIPTION OF THE EVOLUTIONARY MODELS

For the purposes of our study, we calculate a set of evolutionary models for PMS stars with the use of the FRANEC (Frascati Raphson Newton Evolutionary Code) stellar evolution code (Chieffi & Straniero 1989; Degl’Innocenti et al. 2008). Opacity tables are taken from Ferguson et al. (2005) for $\log T[K] < 4.5$ and Iglesias & Rogers (1996) for higher temperatures. The equation of state (EOS) is described in Rogers et al. (1996). Both opacity tables and EOS are calculated for a heavy elements mixture equal to the solar mixture of Asplund et al. (2005). Our models are then completely self-consistent, with a unique chemical composition in the whole structure.

The value of the mixing length parameter, $\alpha$, has been chosen by fitting a solar model, obtained to be $\alpha = 1.9$. Helium and metal contents of the models are those best representative for the LMC, i.e., $Y = 0.25$ and $Z = 0.01$, respectively. The evolutionary tracks have been calculated for a range of masses between 0.2 and 6.5 $M_\odot$. We have interpolated a dense grid
of isochrones from the evolutionary tracks, which covers ages from 0.5 to 200 Myr, with steps of 0.1 Myr below 10 Myr, with steps of 0.5 Myr between 10 and 20 Myr, with steps of 1 Myr between 20 and 50 Myr isochrones, and with steps of 5 Myr above 50 Myr. For ages $\tau < 1.6$ Myr the isochrones do not reach the main sequence (MS) for the highest masses in our sample, while at $\tau \sim 50$ Myr the models cover also post-MS evolutionary phases for intermediate masses.

The evolution of light elements (deuterium, lithium, beryllium, and boron) is explicitly included in this version of FRANEC code. Deuterium-burning plays an important role in the very early stages of PMS evolution for stars with $M \lesssim 0.6$ $M_\odot$. For these stars, the energy provided by the $\text{D}(p, \gamma)^3\text{He}$ reaction is able to slow down the contraction in the initial part of the Hayashi track. This can be seen in Figure 1, where it can be seen that the low-mass end of the PMS isochrones is bent toward higher luminosities for ages younger than $\sim 5$ Myr.

In order to explore the dependence of our findings on the assumed models, consideration apart from the assumed models, we convert them from the physical units ($T_{\text{eff}}, \log g, L_{\text{bol}}$) into absolute magnitudes in the HST/ACS photometric system, by means of synthetic photometry on a grid of atmosphere models. As discussed thoroughly in Paper I, the surface gravity dependence in the computation of optical colors for PMS stars must be accounted for, especially in the low-mass regime. Since atmosphere models are generally expressed in terms of $T_{\text{eff}}$ and $\log g$, synthetic photometry provides a precise consideration of the color dependence on surface gravity, allowing to isolate one unique spectrum as the best representative of the stellar parameters for each point in a theoretical isochrone.

The choice of a reliable atmosphere grid is critical for cool stars (Da Rio et al. 2010), whose spectral energy distributions (SEDs) are dominated by broad molecular bands. Moreover, these spectral features are gravity dependent, and since the stellar surface gravity varies during PMS contraction, optical colors are age dependent for young stars. In Paper I for the derivation of the IMF of LH 95, we utilized the NextGen (Hauschildt et al. 1999) synthetic spectra to convert the theoretical Siess models into colors and magnitudes, extended with the Kurucz (1993) grid for the highest temperatures ($T_{\text{eff}} > 8000$ K). While our ACS photometry of LH 95 in Paper I includes PMS stars down to $\sim 0.2$ $M_\odot$, the detection incompleteness could not be corrected for masses below $\sim 0.4$ $M_\odot$. Therefore, the IMF derived in Paper I was constructed for masses down to $\sim 0.4$ $M_\odot$.

However, in the present study on the age of LH 95 and the age distribution of its low-mass PMS stars, the whole sample of detected stars down to $\sim 0.2$ $M_\odot$ can be used. Therefore, we want to refine the atmosphere models we adopt to compute intrinsic colors for the very low stellar masses.

Baraffe et al. (1998) show that NextGen models cannot reproduce empirical solar-metallicity $V$ versus $(V-I)$ CMDs for MS populations. The problems set in for $(V-I) > 2$, which corresponds to $T_{\text{eff}} \lesssim 3600$ K, and can be attributed to shortcomings in the treatment of molecules. In particular, Allard et al. (2000) illustrate that the main reason for the failure of NextGen spectra in matching the optical properties of late-type stars is the incompleteness of the opacities of TiO lines and water. In that work, they present new families of synthetic spectra (AMES models) with updated opacities, and discuss the comparison with optical and NIR data, finding a good match with empirical MS dwarf colors from the photometric catalog of Leggett (1992).

Da Rio et al. (2010), in their analysis of a multi-color optical survey of the Orion Nebula Cluster (ONC) down to the hydrogen-burning limit, show that, for young, M-type stars: (1) the optical colors are bluer than those of MS dwarfs at the same temperature, (2) intrinsic colors obtained from synthetic photometry using AMES spectra better reproduce the observed colors, and (3) for late M-type stars the AMES colors are even slightly bluer than the observed sequence, needing an additional correction based on the empirical data. As we show later on, the predicted colors from NextGen models are even bluer than the AMES ones for young stars. Moreover, the correction to the AMES colors proposed in Da Rio et al. (2010) is critical for temperatures below our detection limit. As a consequence, we are confident that for the analysis in this paper, the AMES models are in any case more appropriate than the NextGen, and their still present uncertainties do not play a significant role when used on our LH 05 photometry.

Unfortunately, the AMES library is available only for solar metallicity, and therefore we proceed by performing an approximation that provides the best estimate of the correct magnitudes and colors for M-type stars with metallicity $Z = 0.01$. In conclusion, we perform synthetic photometry on the evolutionary models considered here by applying the technique presented by Da Rio et al. (2010) making use of three families of synthetic spectra: (1) NextGen, for an LMC metallicity of $Z \simeq 0.01$, (2) AMES, for a solar metallicity, and (3) Siess et al. (2000) models, for a subsolar metallicity.
Figure 2. Left panel: CMD showing our FRANEC isochrones for ages of 0.5, 1, 3, 5, 10, 30, and 100 Myr converted into the ACS photometric system using the solar-metallicity atmosphere models from the AMES library (solid line), the NEXTGen models of solar metallicity (dashed lines), and the same of LMC metallicity $Z = 0.01$ (dot-dashed lines). Left panel: representation of the displacements in the CMD of identical points on the evolutionary models caused by changing the three grids of atmosphere models used.

(A color version of this figure is available in the online journal.)

(2) NEXTGen for solar metallicity $Z = 0.02$, and (3) combined NEXTGen for $T_{\text{eff}} > 5000$ K and AMES for $T_{\text{eff}} < 5000$ K for solar metallicity. For $T_{\text{eff}} > 8000$ K, all three used grids are extended using Kurucz (1993) spectra.

Results of our transformations are shown in Figure 2 (left) for a number of FRANEC isochrones with ages spanning from 0.5 to 100 Myr. In this figure, it is shown that while the dependence of colors on the used synthetic spectra is negligible for low color terms that correspond to high $T_{\text{eff}}$, the offset increases for late spectral types and young ages. In particular, the more “correct” AMES atmospheres are up to 0.2 mag redder than the NEXTGen ones for the same metallicity, and a decrease of [M/H] by a factor of two (from the solar to LMC metallicity) leads to a similar offset toward bluer colors.

In order to estimate the magnitudes and colors of AMES-transformed isochrones for the LMC metallicity, we compute the displacement produced by changing the metallicity from $Z = 0.02$ to $Z = 0.01$ with the use of the NEXTGen spectra for every point of every isochrone. These displacements are represented by arrows in Figure 2 (right). They can be considered as a fairly good approximation of the real change in the observed colors and magnitudes in relation to $T_{\text{eff}}$ and log $g$, caused by the decrease of metallicity. Therefore, we apply the same offset, point by point, to isochrones transformed with the use of AMES spectra with solar metallicity, in order to derive their counterpart isochrones for $Z = 0.01$.

Figure 3 shows a sample of FRANEC isochrones transformed to the observable plane with the use of the AMES spectra library extrapolated to the LMC metallicity plotted over the CMD of LH 95. A distance modulus of $m-M = 18.41$ mag and average extinction of $A_V = 0.6$ mag, derived in Paper I, are applied to the isochrones. The PMS stars of LH 95, roughly following the 5 Myr isochrone, exhibit a wide apparent spread that spans from the youngest ages available in the models to the MS.

Figure 3. Stellar population of LH 95 with, overlaid, the FRANEC models converted into observable colors and magnitudes by extrapolating the results of the AMES model atmospheres for $Z = 0.01$.

(A color version of this figure is available in the online journal.)
4. FIELD-SUBTRACTION REFINEMENT

Our observations of the control field 2′ away from LH 95 with setup identical to those of the system allowed us in Paper I to perform a thorough statistical removal of the contamination of the observed population by the stars of the local LMC field, on the verified basis that there are no PMS stars in the LMC field. Indeed, this method provided a sufficiently clean CMD of the PMS cluster members of LH 95 (as shown in Figure 3). However, there is an anomalous presence of a small number of sources in the field-subtracted CMD located between the general PMS population and the MS. These sources can be seen in Figure 3 located in the neighborhood of the 100 and 30 Myr isochrones. In Paper I, we hypothesized the origin of these contaminants as a binary sequence of the field MS population, and we demonstrated that their removal does not affect the shape of the derived IMF.

However, for the present study, which focuses on the identification of any age distribution among PMS stars, the nature of these sources and their membership should be reevaluated. In order to study the spread of the field MS in the CMD of both LH 95 and the control field, we construct the Hess diagram of the complete stellar sample observed in the area of the association before any field subtraction. We bin the stars in the color and magnitude axis of the CMD and smooth the counts with a Gaussian kernel of 0.05 mag in $m_{555} - m_{814}$ and 0.16 mag in $m_{814}$. We repeat this process to construct the Hess diagram of the control field. The constructed Hess diagrams are shown in Figure 4. An important difference between the field population covered in the area of LH 95 and that in the control field is that LH 95 is embedded in the H II region named LHA 120-N 64C (or in short N 64C; Henize 1956) or DEM L 252 (Davies et al. 1976), and therefore the reddening of the background field stars in the area of the system should be somewhat higher than those in the control field. However, a comparison of the Hess diagrams of Figure 4 shows that there is no significant variation of the field population between the two areas, and therefore the appearance of the nebula does not seem to affect significantly the field MS. This suggests that the LMC field population is mainly distributed in foreground with respect to LH 95. Nevertheless, a small fraction of field stars could be located behind the nebula and be affected by a larger extinction.

In order to test this possibility and further eliminate the sample of PMS cluster members from potential field MS contaminants, we consider diagonal cross sections of width 0.33 mag in the direction defined by the line $m_{814} = -1.92 \times (m_{555} - m_{814})$, and we fit with Gaussians the distributions of the stars along each of these strips. A sum of two Gaussian distributions is considered for the fit of the CMD of LH 95 and a single Gaussian for that of the control field. The centers and standard deviations, $\sigma$, of the best-fitting Gaussians are drawn over the Hess diagrams of Figure 4. We compare the $\sigma$ of the best-fitting Gaussians to the MS, as a measurement of the broadening of the MS, between the CMD of LH 95 and that of the control field in Figure 5, where we plot $\sigma$ for each selected strip as a function of the corresponding average luminosity. In this figure it can be seen that $\sigma$ and therefore the MS broadening are systematically smaller in the control field than in LH 95, in agreement with what we discussed above. This difference demonstrates quantitatively that the contamination of the CMD of LH 95 by field stars cannot be 100% removed, and this applies mostly to the reddest part of the field MS in the area of LH 95 due to higher reddening. As a consequence, we further remove from the sample of cluster-member sources those that fall blueward of the minima between the two Gaussian fits of the CMD of LH 95, shown in Figure 4 (left). This second selection border is drawn on the CMD of the area of LH 95 in Figure 6; we prolong this border at lower and higher luminosities, as shown by the dashed line in Figure 6.
tributions (close to the cut shown in Figure 6) is still present.

certainty in the membership assignment on the tail of the dis-

subtraction refinement with respect to that of Paper I, some

analysis.

the field and the system.

for LH 95, based on the results of Gaussian fits to the distribution of the MS in

field subtraction in Paper I. Here, we refine our selection of true stellar members

stars could not be completely removed within the application of the statistical

broadening of the MS in the control field than in LH 95 demonstrates that field

the area of LH 95 (solid line) and the control field (dashed line). The smaller

Figure 6. These sources will be excluded from our subsequent

remnants, plotted with gray dots to the left of the MS border in

handful of stars at 

the bright end we arbitrarily define a cut in order to remove a

source of 22 mag, undoubtedly located on the MS.

There are 162 stars, which we consider to be field-subtraction remnants, plotted with gray dots to the left of the MS border in Figure 6. These sources will be excluded from our subsequent analysis.

We stress that, despite the improvement of this field-subtraction refinement with respect to that of Paper I, some uncertainty in the membership assignment on the tail of the distributions (close to the cut shown in Figure 6) is still present. This effect, which might affect the distribution of the oldest stars, should not be significant. The population, within the central region, considered MS field is about as numerous as that considered part of the young system (1904 and 1942 stars re-

As a consequence, the number of young members erroneously removed on account of being too blue should be roughly compensated by a similar number of MS stars redder than the minima between the two Gaussians (Figure 4). This guarantees that the number of old PMS members of stars is fairly accurate, but could lead to an underestimation of the ages of some of those. In any case, since in this work we demonstrate a significant age spread for LH 95, this fact would just reinforce our findings.

5. AGE DETERMINATION OF PMS STARS

In this section, we apply our new maximum-likelihood technique for the derivation of the ages of the PMS stars of LH 95 from their photometry. This method allows the determination of the age of a PMS star from its position in the CMD in a probabilistic manner, accounting for the physical and observational effects that affect the colors and magnitudes of all observed PMS stars.

In order to estimate the age of a young stellar system when only photometry is available, it is common practice to simply overlay a set of isochrones to the CMD and, by interpolation, assign an age to all the members. This approach, however, can lead to inaccurate results (Gouliermis et al. 2010). This is because every star is scattered from its original position in the CMD by a number of factors, and specifically by photometric errors, differential reddening, unresolved binarity, stellar variability, non-photospheric excess due to accretion, and scattered light from circumstellar material.

These effects, when not accounted for, bias the observed age distribution, both increasing the observed spread and systematically shifting the mean cluster age. These factors can be disentangled only on the basis of individual stars, which is the case when spectroscopic measurements are available for all stars. However, such observations on large numbers of stars are not technically feasible, and photometric observations are the most appropriate for collecting data on vast stellar numbers, like in the case of LH 95. In this case, a convenient way to estimate the overall effect of observational and physical biases to the original CMD positions of PMS stars is to compare the observed data with synthetic CMDs, which are constructed by the application of these biases to one or more theoretical PMS isochrones. There are some examples of such an approach presented in the literature, especially for the investigation of the relation between the observed luminosity spread and the true age spread among PMS star members of young stellar systems (e.g., Burningham et al. 2005; Jeffries et al. 2007; Hennekemper et al. 2008; Hillenbrand et al. 2008; Weidner et al. 2009).

In this context, Naylor & Jeffries (2006, hereafter NJ06) presented a rigorous approach to derive the age of young stellar systems accounting for photometric errors and binarity. They present a maximum-likelihood method for fitting two-dimensional model distributions to observed stellar CMDs. In their application, the two-dimensional models are probability density distributions in the CMD, obtained by applying bina-

rity to a PMS isochrone of a given age, considering that the model representing a coeval population in the CMD is not a line—like a theoretical isochrone—but a two-dimensional
density distribution. There are several advantages in the application of this approach. First, it enables one to account rigorously for the photometric errors in both color and magnitude when the fit is performed. Second, binarity is considered in the fitting process itself. Finally, this method allows the evaluation of its goodness of fit.

In the present study, we further develop the method of NJ06, by extending its application to a realistic case of a PMS population, including in the computation of the two-dimensional distribution of this approach. First, it enables one to account rigorously for the photometric errors in both color and magnitude when the fit is performed. Second, binarity is considered in the fitting process itself. Finally, this method allows the evaluation of its goodness of fit.

5.1. From Isochrones to Density Distributions in the CMD

As in NJ06, we wish to construct a stellar density distribution in the CMD for every age in the considered set of evolutionary models. This distribution represents the probability for a star of an assumed age to be found at a given point in the CMD. We consider the two families of evolutionary models (FRANEC and Siess) that we converted into magnitudes in the HST/ACS photometric system in Section 3.1. For every considered age, we populate the appropriate isochrone with $5 \times 10^5$ simulated stars. We assume the IMF of LH 95 derived in Paper I and correct it for binarity. The change in the measured IMF slopes caused by binarity was already investigated in Paper I, through a Monte Carlo simulation. Here, we adopt these results, summarized in Table 1. We found that the IMF for intermediate and high-mass stars ($M > 1M_{\odot}$) is almost insensitive on binaries, whereas for low-mass stars, the slope changes of up to 0.4 units adding an increasing fraction of unresolved binaries. Neglecting binaries leads to a shallower measured IMF.

For smaller masses, not having a direct measurement of the IMF in LH 95, we assume a value $x = 0.3$ compatible with the Kroupa (2001) IMF. This value was also assumed in Paper I to compute the binarity-corrected IMF slopes reported in Table 1. It is worth noting that small changes in the slope of the assumed IMF do not influence significantly the ages for single stars derived using our method (see Section 5.2), but could affect the evaluation of the goodness of fit (Section 6) we use to determine the age spread.

After simulating a population of stars for every age (for the two families of evolutionary models and all the binary fractions $f$ reported in Table 1), we displace with a random process the stars in the CMD as dictated by several factors. They are discussed below.

### 1. Rotational variability

We simulate statistically the variation of magnitudes in the F555W and F814W bands due to variability caused by stellar rotation. We consider the results of Herbst et al. (2002). These authors surveyed the PMS population of the ONC in the $I$ band over 45 nights, obtaining the light curves for about 1500 candidate members. We consider their catalog of peak-to-peak variations. A single imaging of a young stellar system, like LH 95, statistically measures the flux of variable stars in random phases of their variation. Considering, at this point, variability only due to dark spots on the stellar surface, the measured fluxes are equal or fainter than those predicted if no variability is considered. Therefore, if we assume for simplicity sinusoidal light curves, the statistical distribution of $I$-band variations for a star with a measure peak-to-peak $\Delta I_{\text{max}}$ is the distribution of $\Delta I_{\text{max}} / \sin \alpha$ with $0 < \alpha < 2\pi$. We compute this for every star of the catalog of Herbst et al. (2002) and combine the derived distributions. The result, as shown in Figure 7(a), is the observed $I$-band variation distribution that we assume valid for single epoch photometry in a young stellar system. The peak at very low $\Delta I$ is mainly due to the large fraction of stars not showing significant rotational variability. We displace the brightness in the F814W filter of each of the simulated stars along each isochrone by a value randomly drawn from this distribution.

The amplitude of the $V$-band variation due to dark spots is generally larger than in $I$, and the ratio between the two depends, in general, on the temperature of the starspot, on the temperature of the photosphere, and the filling factor of the spot. According to Herbst et al. (1994), the ratio $\alpha = \Delta I / \Delta V$ spans from 0.39 to 0.88. The variations in the two bands should be correlated if the observations were close to co-temporal. In our photometry, however, as mentioned in Section 2, $V$ and $I$ observations were taken separately, one 4 days after the other. According to Herbst et al. (2002), the rotational period distribution of T Tauri stars is bimodal, with a peak at $P \sim 2$ days and another at $P \sim 8$ days. Even considering the longest periods, the time delay of our observations covers about half period. As a consequence, we cannot consider the $V$- and $I$-band variations to be correlated. We therefore simulated the $V$-band variation distribution from the $I$-band one, multiplying each of the simulated $I$ values for a factor $1/\alpha$, where $\alpha$ was randomly chosen within the range 0.39–0.88. Then these $V$ variations were applied randomly (regardless of $\Delta I$) to all the simulated stars.

### 2. Accretion

Ongoing accretion from the circumstellar disk onto the stars also affects the observed fluxes. In the optical wavelength range, two components contribute to the accretion excess: an optically thin emission generated in the infalling flow, and an optically thick emission from the heated photosphere below the shock (Calvet & Gullbring 1998). The optically thin emission is characterized by a line-dominated spectrum, similar to that of an H II region; in particular the bright Balmer emission is commonly used to derive mass accretion rates for PMS stars (e.g., De Marchi et al. 2010). In $V$ and $I$ band, the optically thick emission dominates, producing an excess, in spectroscopy referred to as “veiling,” analogous to the presence of a hot spot on the stellar photosphere. For a given accreting star, the spectrum and intensity of the excesses may vary with time, due to both variations in the mass accretion, and stellar rotation.

In order to model the overall effect of accretion on a large sample of low-mass stars, we utilize the results of Da Rio et al. (2010). In this work, the stellar population of the ONC was studied from multi-band optical photometry and spectroscopy. In particular, for every star, the veiling

### Table 1

| f     | $x\ (M < 1M_{\odot})$ | $x\ (M \geq 1M_{\odot})$ |
|-------|----------------------|-------------------------|
| 0.0a  | 1.05                 | 2.05                    |
| 0.2   | 1.20                 | 2.05                    |
| 0.4   | 1.30                 | 2.05                    |
| 0.5   | 1.35                 | 2.1                     |
| 0.6   | 1.40                 | 2.1                     |
| 0.8   | 1.45                 | 2.1                     |

**Notes.** In the units we use to express IMF slopes, a Salpeter (1955) IMF has a slope $x = 1.35$.

a Measured IMF from Paper I.
excesses were determined from the comparison of the observed colors with the intrinsic ones, after modeling the color displacements caused by accretion and extinction. For the ONC, the authors find about 60% of the members showing evidences of accretion veiling. We consider this stellar sample, and restrict to the mass range $0 < M \leq 1$ (Woitas et al. 2001), and do not allow binary systems to form with mass ratio $q < 0.2$. We let the binary fraction be a free parameter, and compute different sets of two-dimensional models for different choices of its value. This is because, as we discuss later in detail, the assumed binary fraction influences significantly the derived ages, and therefore we present and compare findings obtained with different choices of this parameter. We assume a unique binary fraction and mass ratio distribution for the whole PMS mass range. In principle, high-mass stars are known to have both higher binary fraction and higher binary mass ratio than intermediate and low-mass stars (e.g., Lada 2006; Zinnecker & Yorke 2007). However, this does not affect our treatment, since in our analysis of stellar ages we only consider PMS stars and not the “evolved” high-mass MS stars of LH 95.

4. Differential extinction. In Paper I, we derived the extinction distribution for LH 95, measuring the reddening of the UMS from the zero-age main sequence (ZAMS). In that work not all the UMSs were considered to derive the $A_V$ distribution, but some of the reddest were removed, since we suspected them to be Herbig Ae/Be stars. The average extinction is quite low, equal to $A_{555} \sim 0.6$ mag, and its distribution spans from $A_{555} = 0$ to $A_{555} \sim 1$. We use the relation $A_{555}/A_{814} = 1.85$, as it was computed in Paper I for the ACS photometric system, using the extinction law of Cardelli et al. (1989). In our simulations, we apply a random value of extinction to each star, drawn from the real measured $A_{555}$ distribution of LH 95.

5. Source confusion. LH 95 is a loose association, not centrally concentrated, which comprises at least three subclusters of somewhat higher stellar densities. As a consequence, source confusion due to crowding does not play a dominant role. Nevertheless, we include in our simulations a 5% source confusion due to crowding, allowing for random confusion also between binaries and between binaries and isolated stars. We stress that, since this additional pairing is done randomly, because of the IMF the modeled source confusion is also magnitude dependent, favoring the disappearance of faint systems.

For the computation of the two-dimensional stellar densities, we do not include photometric errors as additional sources of spread. While the aforementioned biases, which we include in the simulations, can be considered as part of the model, the photometric errors are part of the data, which the model is fitted to, and therefore photometric uncertainty will be consider later in the fitting process (Section 5.2).

It should be finally noted that the constructed two-dimensional distributions are not re-normalized to, e.g., a total integral of 1, because isochrones of different ages have different...
maximum mass. Instead, we use the same sample of five million simulated stars with masses drawn from the assumed IMF and we consider only the fraction included within the limits of each isochrone. In this manner, although the total number of synthetic stars changes slightly among different two-dimensional distributions, the number of stars per mass interval is constant. In other words, since our method to assign stellar ages and determine the cluster age is based on relative comparison of probabilities, an overall normalization of all the models is completely unnecessary.

In Figure 8, we show an example of the result of the application of our simulations to a 2 Myr isochrone from the Siess family of models. In this figure, it is evident that the CMD spread of a “perfect” coeval PMS stellar population, produced by the biases discussed above, can generate an observed dislocation of the stars of more than 1 mag in luminosity and up to 1 mag in color. The bifurcation seen at the faint end of the simulated distribution is an effect of binarity. This is due to the fact that the evolutionary models do not cover masses lower than 0.1 $M_\odot$. As a consequence, our random selection of binary pairs is unable to simulate binaries with the second component below this limit. This effect does not alter the results of our analysis, because, as it is seen in Figure 8, our photometry does not reach these masses.

### 5.2. Derivation of Stellar Ages: The Single Star Case

In this section, we use statistics to address the issue of deriving the most likely age of a single PMS star, for which the only information is its measured brightness and the corresponding photometric errors. Specifically, we derive the age of each star in our catalog using a maximum-likelihood method by comparing the observed position of each star in the CMD with a complete set of simulated two-dimensional stellar distributions, constructed from theoretical isochrones with the method described above. Each of these simulated distributions provides an estimate of the probability of finding a star in a specific CMD location, providing that we know its age. In Section 5.4, we will then generalize this approach to derive the best age of the entire cluster, considering all the measured stars simultaneously.

In practice, if the photometric errors were zero, the most likely age would be that of the modeled distribution that gives the highest probability in the CMD position of the star. In reality, since there are errors in the measured $m_{B55}$ and $m_{B14}$ magnitudes of the star, the observed position of the star in the CMD is also probabilistic, defined by the Gaussian distributions of the photometric errors. Therefore, if we consider a star observed in the position $(V_1, I_1)$ on the CMD and a model stellar distribution of a given age $t$, for every point $(V_j, I_j)$ of the CMD space there are two probabilities that should be specified: (1) the probability, defined by the two-dimensional model, that a star of age $t$ is located in $(V_1, I_1)$, and (2) the probability, defined by the Gaussian distributions of the photometric errors, that the star we observe in $(V_1, I_1)$ is instead located in $(V_j, I_j)$ but displaced by the photometric errors. The product of these two probabilities, integrated over the entire CMD space, provides the likelihood that the considered star has an age equal to $t$. This likelihood, normalized on all the possible ages, represents the exact probability that the star has an age between $t$ and $t + dt$. Since the photometric errors in the color–magnitude space are correlated, we apply this method in the magnitude–magnitude plane. Mathematically, we define the likelihood $L_i(t)$ for the $i$th star in our sample as function of the age $t$ as

$$L_i(t) = \int_{V_1} V_j \rho(V, I, t) U_i(V - V_i, I - I_i) dV dI,$$  \hspace{1cm} (1)

where $\rho(V, I, t)$ is the surface density computed from the isochrone of age $t$, and

$$U_i(V - V_i, I - I_i) = \frac{1}{2\pi \sigma_V \sigma_I} e^{-\left(\frac{(V - V_i)^2}{2\sigma_V^2} + \frac{(I - I_i)^2}{2\sigma_I^2}\right)}$$  \hspace{1cm} (2)

is the two-dimensional Gaussian representing the photometric error of the $i$th star.

In Figure 9, we show an example of the likelihood function $L_i(t)$ computed numerically for three stars in the LH 95 catalog, assuming two-dimensional models obtained from the FRANEC grid and assuming a binary fraction $f = 0.5$. The age assignment for every star is performed by maximizing $L_i(t)$. For example, the three stars whose $L_i(t)$ functions are shown in Figure 9, will have most probable ages of 1.3 Myr, 5.5 Myr, and 19 Myr.

In order to estimate the accuracy in the age determinations, we isolate for every likelihood function the range, in log $(t)$ that includes 68% of the area under the curve, in analogy with the $\pm 1\sigma$ interval for a normal distribution. Since for some stars the shape of $L_i(t)$ is highly skewed, we choose to define the limits of the interval to delimit identical areas of 34% of the total to the left and to the right from the peak of the distribution. We then call $\epsilon$ the width of the range, in units of log $(t)$. Indicatively, if a $L_i(t)$ were normal, $\sigma(t)$ would be equal to $\epsilon/2$.

Figure 10 (left) shows the derived best-fit stellar ages in the CMD, assuming the evolutionary models of FRANEC and a
Figure 9. Example of likelihood functions, $L_i(t)$, for three stars from our photometric sample. The maxima of the functions indicate the assigned best-fit ages for the sources, equal to 1.3, 5.5, and 19 Myr, respectively, in this example.

binary fraction $f = 0.5$. In Figure 10 (right), we show the corresponding variation of $\epsilon$ along the CMD. It is evident that, whereas the age assignment turns out to be relatively precise in the PMS regime, for higher masses, close to the PMS–MS transition ($m_{814} \lesssim 21$) the uncertainty increases up to $\epsilon \approx 1.5$ dex. This is due to the fact that PMS isochrones of different ages are closer to each other in this range, and tend to overlap in the MS regime. The higher values of $\epsilon$ on the blue side of the sequence are also due to the fact that theoretical isochrones of ages $\gtrsim 10$ are also close to each other, since PMS evolution is significantly slower for these ages than for younger stages of PMS contraction.

We compare the distributions of the best-fit ages for our photometric sample changing the binary fraction and the original family of evolutionary models. This is shown in Figure 11, where we also indicate the average age for each distribution. It this figure is evident that increasing the binary fraction leads to the prediction of older ages. This confirms the findings of Naylor (2009). Indeed, a higher fraction of binary systems produces a simulated two-dimensional distribution shifted toward higher luminosities, hence older measured ages. Passing from $f = 0$ to $f = 0.8$, we predict a difference in age of up to 0.2 dex (which is 60% older ages.).

In the same figure, we also overlay the distribution of ages computed only by interpolation between theoretical isochrones, without the addition of any observational and physical bias except for the average extinction of LH 95. In this case, we find that the average age is higher than that derived from our statistical method assuming $f = 0$. This is due to the effect of dark spots and accretion. The first biases the predicted luminosities toward fainter values and the second toward brighter luminosities and bluer colors. In particular, the color effect caused by veiling affects more the position in the CMD than the increase in the fluxes, requiring an older isochrone to fit the observed sequence.

In conclusion, we find that, in the derivation of stellar ages for a young PMS cluster, neglecting binaries leads to an underestimation of the ages, neglecting variability and accretion to an overestimation of those. The two effects counterbalance, in our case, for a binary fraction $f = 0.5$ and assuming starspot and accretion values from the 2 Myr ONC.

5.3. Spatial Variability of Stellar Ages

As we mentioned in Section 2, LH 95 is not a centrally concentrated cluster. It presents three main concentrations of PMS stars, named in our previous work “subcluster A, B, C,” located, respectively, at the eastern, western, and southern edges of LH 95. They include about half of the stellar population of the region, while the rest of the population is distributed over the area between the three. Since the projected distance between the stellar clumps is quite large (15–20 pc), it could

Figure 10. Left: color–magnitude diagram with the PMS stars color-coded according to their individual ages derived with our maximum-likelihood method. The model distributions were derived with the use of theoretical isochrones from the FRANEC family of models. The assumed binary fraction is $f = 0.5$. Right: the same CMD with the stars color-coded according to their $\epsilon$ parameter (see Section 5.2).
Figure 11. Distribution of maximum-likelihood stellar ages for the LH 95 members, computed using two-dimensional CMD distribution models, based on FRANEC (upper panel) and Siess (lower panel) evolutionary models. We assume different binary fraction values $f$ from 0.0 to 0.8. For each family of isochrones, we also show the distribution of ages computed only by interpolation of isochrones in the CMD, without accounting for physical and observational biases (black lines). The vertical lines (colored appropriately) in each panel indicate the mean value of the age for each plotted distribution.

Figure 12. Comparison between the age distribution of the entire region and the three denser subclusters, for FRANEC evolutionary models and $f = 0.5$.

be plausible that star formation did not occur simultaneously all over the entire region, but sequentially. A Galactic example of such scenario can be found in the Orion complex (Briceno 2008), where different populations, from 2 to 10 Myr old, are distributed over a similar range of distances.

We investigate if the stellar sub-groups in LH 95 present different average ages, in order to rule out the possibility that the measured broad age distribution for the entire region (Figure 11) is simply due to a superposition of populations with different ages, instead of the effect of observational biases and, maybe, a real age spread. To this purpose, we isolate the members in the three subclusters, using the same selections as in Paper I. Then we consider the distribution of ages limited to each of the three subregions, in comparison with that of the entire LH 95. This is illustrated in Figure 12. It is evident that, in each region, the distributions are very similar, spanning from less than 1 Myr to 25 Myr, with a peak at $\sim$ 5 Myr. The only evident anomaly is an overabundance of young members for subcluster B. However, this group of stars encloses the densest concentration of PMS stars in the whole region, with more than 150 stars within $\sim$ 2 pc. Because of this, stellar confusion due to crowding should be particularly dominant in this area, producing a number of overluminous stars due to unresolved superposition, and, thus, a fraction of apparently younger members. This is, however, not a problem for the analysis of the cluster age (Section 5.4).

In our CMD modeling (Section 5.1), we have purposefully included the effect of crowding affecting 5% of members. The overabundance of young stars in subcluster B is caused by about 40 stars, which are $\sim$ 2% of the total stellar sample.

Considering this, we exclude the presence of any spatial variations of the ages in LH 95. Despite the irregular density of stars and the relatively large size of the region, LH 95 must have formed as a whole, in a unique episode of star formation.

5.4. Age of the Entire Cluster

In Section 5.2, we derived the most likely age for every star in the sample. Here, we derive the best-fit age for the entire stellar system. In a similar way as in NJ06, we compute this age by maximizing the “global” likelihood function $L(t) = \prod L_i(t)$. To compute $L(t)$, we restrict our treatment to the stellar sample with $22 \leq m_{555} \leq 27$. In particular, we want to eliminate any biases introduced by detection incompleteness, which would lead to the detection at the faint end of more young bright members than old faint ones. The detection completeness reaches $\sim$ 90% at $m_{555} = 27$ (Paper I), and we consider this threshold as a reasonable limit for our analysis here. We stress that, as in Naylor (2009), we could include the detection completeness function in the two-dimensional models derived in Section 5.1, and then recover the implied cluster properties using our maximum-likelihood procedure considering all the stars. In this way, we could test for bias also as a function of our chosen completeness limit. Nevertheless, considering the large stellar sample we have at disposal, we prefer to limit our analysis to the luminosity range where the detection incompleteness does not affect our photometry. In this way, we remove any additional uncertainty in the cluster age caused by a possible inaccuracy of the measured completeness function.

In addition, we do not include in the analysis the bright members of both the MS and the PMS–MS transition, for two reasons: (1) as shown in Figure 10 (right), the accuracy of the age estimation for these sources is poor, and (2) these sources...
of up to 60% with increasing binarity (depending on the considered isochrone models). The peak values of $\log L(t)$ given in Table 2 do not have an absolute meaning, given the arbitrariness of the normalization of the two-dimensional density models $\rho(V, I, t)$.

In Figure 14, we show again the observed CMD, with the best-fit two-dimensional density models from FRANEC (left) and Siess (right) grids of isochrones superimposed for $f = 0.5$, for the luminosity range we considered in our analysis. The two models are very similar, with FRANEC producing a slightly overdensity at the reddest end of the considered magnitude range.

From Figure 14, it is evident that an additional spread, besides that modeled in the two-dimensional density distribution, is evident in our data. In particular, there are observed members at the sides of the sequence which may not be well modeled by the simulated distributions for a single cluster age. This may be indicating of an additional age spread in LH 95. In the next sections we study this possibility.

6. AGE SPREAD OF PMS STARS IN LH 95

6.1. Verification of an Age Spread

In the previous section, we have determined the average age of LH 95 isolating the two-dimensional model that provides the best fit to the observed sequence, for two families of theoretical evolutionary models. In this section, we wish to investigate whether the spread of the PMS stars in the observed CMD is wider or not than that of the simulated two-dimensional density distributions as they are produced by accounting for the sources of CMD broadening discussed in Section 5.1. after we additionally apply the photometric uncertainties. If indeed the observed CMD broadening of PMS stars is wider than the synthetic derived from a single isochrone, under the hypothesis that we do not have underestimated the sources of apparent luminosity spread in the CMD, this would suggest the presence of a real age spread in LH 95. The method applied in this section, described below, also allows us to qualify the goodness of our fit.

This method can be compared to a standard $\chi^2$ minimization, in which the best-fit model is obtained by minimizing the $\chi^2$, and this value is then compared to the expected distribution of $\chi^2$ given the degrees of freedom. A measured $\chi^2$ which significantly exceeds the expected distribution indicates a poor fit, while a value too low suggests that the fit is suspiciously good, which is the case, e.g., of overestimated errors. In our statistical framework, the procedure is qualitatively similar, with two main differences: (1) since the measured $\log L(t)$ is maximized, a poor fit is found when the measured value is lower than expected and (2) the expected distribution of $\log L(t_{\text{best-fit}})$ is not analytical but can be derived numerically. In general,
Figure 14. Observed CMD with the two-dimensional density models that provide the best fit to the observed sequence of PMS stars, assuming binarity of $f = 0.5$ for the FRANEC (left panel, age 3.9 Myr) and the Siess (right panel, age 3.8 Myr) family of evolutionary models. The dashed lines denote the subsample of the photometry used in this analysis. For clarity, only the portion of the two-dimensional models within this luminosity range is shown.

this distribution varies for different assumed two-dimensional models, and depends on the number of data points.

We compute the expected distribution of log $L$ as follows. We consider the two best-fit two-dimensional models shown in Figure 14, limiting ourselves, for the moment, to the case of an assumed binary fraction $f = 0.5$. We limit the model to the same range used to derive the cluster age ($22 \leq m_{555} \leq 27$), and simulate a population of stars randomly drawn from the distribution of the model itself, as numerous as the LH 95 members in the same luminosity range. We consider each of the simulated stars, and displace its position in the CMD by applying the average photometric error measured in its immediate CMD area. In particular, for each simulated source we consider the 60 observed stars closer to it in the magnitude–magnitude plane. Then we randomly assign the photometric error of one of these 60, with a weight proportional to the inverse of the distance in magnitudes. We then apply the same method described in Section 5.4 to the simulated population, deriving log $L$. We iterate this procedure 100 times, each time randomly simulating a different population of test stars from the same best-fit two-dimensional model. In this way, we obtain 100 values of log $L$ describing the expected distribution of this variable for the considered model, which is compared to the measured value computed for the observed sequence of PMS stars.

For the two models shown in Figure 14, i.e., the FRANEC 3.9 Myr and the Siess 3.8 Myr models for $f = 0.5$, we obtain predicted log $L$ values of $2739 \pm 20$ and $2473 \pm 19$, respectively ($\pm 1\sigma$ interval). In comparison, the measured values of $2533$ and $2416$ (see Table 2) are significantly lower, at $\sim -10\sigma$ and $-3\sigma$ from the expected values. This implies that the observed sequence cannot originate from the best-fit models by additionally applying photometric errors, but that the stars tend to be statistically located in regions of the CMD where the two-dimensional model has lower density $\rho$ (see Equation (1)), and hence a lower average log $L_i$ for the stars. For guidance, the low-density regions of the two-dimensional models in Figure 14 are those indicated with lighter blue colors, i.e., at high luminosity, where the IMF predicts fewer members, and at both sides along the two-dimensional distribution in the CMD, at large distance from the peaks of the two-dimensional sequence itself. Since the IMF we use to compute the two-dimensional models has been derived from the actual observed data, we exclude the possibility that the low measured value of log $L$ is due to an overabundance of intermediate-mass stars in comparison to low-mass stars. As a consequence, our results suggest that the observed sequence of PMS stars in the CMD is broader than that predicted by the two-dimensional model for a single age. There are two possible interpretations for this: (1) LH 95 is formed by a coeval population, and the extra spread in the observed CMD is due to additional sources of apparent broadening in the CMD besides those we have included in our models and (2) LH 95 hosts a real age spread among its PMS stars. From hereafter we will assume the second possibility, and, under the validity of our modeling, we determine the age spread of the population.

6.2. Evaluation of the Age Spread

In order to quantify the age spread, we include in our construction of the two-dimensional model distributions, apart from the sources of CMD broadening of PMS stars previously considered (Section 5.1), also a distribution in ages. Specifically, we consider the best-fit age previously measured for the cluster (Table 2) along with a Gaussian distribution of ages peaked at the best-fit age with standard deviations varying from $\sigma_{age} = 0.1$ Myr to $\sigma_{age} = 3$ Myr. We construct, thus, broader
two-dimensional model distributions by essentially adding together different single-age models, computed in Section 5.1, each weighted according to the Gaussian distribution. For each value of $\sigma_{\text{age}}$, we compute the log $L$ from our data on the newly constructed two-dimensional model, as well as the expected distribution of log $L$, from simulated stars drawn from the model, as described above. We stress that the expected distribution of log $L$ is not unique, but varies for each assumed value of $\sigma_{\text{age}}$, being derived from the two-dimensional model computed for the exact value of $\sigma_{\text{age}}$.

For every value of $\sigma_{\text{age}}$, we compare the “observed” log $L$ (i.e., derived from the data) with the predicted one. As mentioned earlier, if the first is significantly smaller than the second, it indicates a poor fit; vice versa, a significantly larger value implies a suspiciously good fit, analogous to a $\chi^2 \ll 1$. The standard deviation of the predicted log $L$ distribution allows us to quantify these differences in a meaningful way. This is simplified by the fact the predicted log $L$ is normally distributed, due to the central limit theorem, since log $L$ is the sum over a large sample of stars of the individual single-star likelihoods log $L_i$.

In Figure 15, we present the difference between the observed and predicted log $L$, normalized on the width of the predicted distribution, as a function of the assumed age spread $\sigma_{\text{age}}$, for the two families of evolutionary models and assuming $f = 0.5$. The values of $\sigma_{\text{age}}$ for which the difference in the y-axis is zero provide the best-fit value of the age spread, the and the projection on the x-axis of the ±1$\sigma$ (dotted lines in the figure), the associated uncertainty. We find that, for FRANEC models, the measured age spread is $\sigma_{\text{age}} = 1.8$ Myr (FWHM ± 4.2 Myr), while for Siess models this decreases to $\sigma_{\text{age}} = 1.2$ Myr (FWHM ± 2.8 Myr). In both cases, the uncertainty in $\sigma_{\text{age}}$ is about 0.2 Myr.

The observed CMD, with the best-fit two-dimensional density models from FRANEC (left) and Siess (right) grids of isochrones superimposed is shown again for $f = 0.5$ in Figure 16, but with the best-fit two-dimensional models constructed by including the measured age spread of 0.5 dex. In this figure is evident how the model density probabilities now follow quite nicely the observed broadening of the PMS population of LH 95, within the magnitude range $22 \leq m_{655} \leq 27$, which is considered for the evaluation of the age and age spread of the system. The difference of these plots with those of Figure 14, where the best-fit two-dimensional models were constructed assuming coeval PMS populations, is significant. By adding an age spread, the model distributions in the CMD now reproduce well the observed sequence.

Our method predicts a 50% higher age spread assuming FRANEC models than assuming Siess models, and the difference is larger than the uncertainties of each of two values. This is a consequence of the particular shape of the FRANEC young isochrones for the lowest masses, where the modeling of the deuterium burning leads to a brighter predicted luminosity. This can be seen already in the original isochrones in the H-R diagram (Figure 1), and is responsible for the overdensity seen in both Figures 14(a) and 16(a) at $V - I \sim 3$ and $I \sim 24$. In this region of the CMD, the models predict more stars than actually observed, regardless the assumed age spread. Since the observed population statistically populates more the positions of the CMD where the density of the two-dimensional models is lower, the measured likelihood turns out to be smaller than the predicted one even for the “right” age spread. This leads to an overestimation of the value of $\sigma_{\text{age}}$.

We remark that this is a general behavior of the method we have proposed: our technique may overestimate the overall age spread. This occurs in the case where the isochrones do not follow the observed sequence, but have a systematically different shape with respect to it. In fact, the condition we imposed to constrain the measured $\sigma_{\text{age}}$ is based on the requirement that our data and a simulated population of test stars obtained from the best two-dimensional density model (including an age spread) produce similar values of “global” likelihood log $L$. If the two-dimensional models have a shape which strongly differs from the observed sequence, the model would never reproduce the data, but there will be always an assumed value of age spread which leads to similar values of likelihood as for the observed sample. In this case, the derived age spread might be higher than the real one. As we discuss in the next section, we can quantify any systematic differences in the shape of the isochrones with respect to the data, and estimate the correct age spread, by analyzing the age properties of the population in smaller luminosity ranges.

We investigated how the assumed binary frequency $f$ affects the derived age spread. We have assumed $f = 0.4$ and $f = 0.6$ and reapplied the method described above. We stress that, although we previously described our results for a larger range of assumed binary ratios, in the mass range relevant for our study, the value of $f$ is not expected to differ much from 50% (Lada 2006). The results are presented in Table 3. We find that a higher fraction of binaries leads to a larger derived age spread and, in contrast, a lower $f$ to a shorter $\sigma_{\text{age}}$. Both the results are a consequence of the dependence of the average cluster age as a function of $f$; given that the distance between consecutive isochrones decreases with age. A given spread in the CMD results in a larger age spread if the central isochrone is older, and vice versa.

6.3. Luminosity Dependence of the Age and Age Spread

We want to verify the presence of systematic changes of both age and age spread with luminosity. This effect, which can be caused by uncertainties in the theoretical evolutionary models, as well as in our modeling of the broadening effects in the CMD, can lead to an overestimation of the measured age spread. To demonstrate this dependence, we divide the photometric sample in the considered magnitude range of
A family of two-dimensional models with different assumed photometry used in the analysis.

For an assumed binarity of $f = 0.5$, the model obtained from FRANEC isochrones, with an average age of 3.8 Myr and a Gaussian age spread of $\sigma_{\text{age}} = 0.18$ Myr. Right: the model obtained from Siess isochrones, with an average age of 3.9 Myr and $\sigma_{\text{age}} = 0.12$ Myr. The dashed lines denote again the subsample of the photometry used in the analysis.

### Table 3

| $f$ | FRANEC Models | Siess Models |
|-----|---------------|--------------|
|     | $\text{Age}$ (Myr) | $\sigma_{\text{age}}$ (Myr) | $\text{Age}$ (Myr) | $\sigma_{\text{age}}$ (Myr) |
| $f = 0.4$ | 3.6 | 1.4 ± 0.2 | 3.6 | 1.0 ± 0.2 |
| $f = 0.5$ | 3.8 | 1.8 ± 0.2 | 3.9 | 1.2 ± 0.2 |
| $f = 0.6$ | 4.1 | 2.3 ± 0.25 | 4.0 | 1.4 ± 0.2 |

$22 \leq m_{555} \leq 27$ into five sub-ranges 1 mag wide each, and we repeat the method described in Section 5.4 to derive the best-fit age within each subsample. We then apply the method of Section 6 to derive the corresponding age spreads, simulating a family of two-dimensional models with different assumed $\sigma_{\text{age}}$ around the average age measured for each magnitude interval.

The results are given in Table 4, again for both the FRANEC and Siess families of evolutionary models for an assumed binary fraction $f = 0.5$. We find a clear variation of both average age and measured age spread with luminosity. In particular, for the bright members, close to the PMS–MS transition, which correspond to masses in the range $1 \lesssim M/M_\odot \lesssim 2$, the derived ages are systematically younger than those for fainter luminosities (and lower masses). The differences in age are of the order of 4 Myr. Also, the measured age spread tends to be significantly larger than the average one for bright stars, and smaller for the lowest masses. In particular, for $22 \leq m_{555} \leq 23$ the age spread derived with our method is larger than $\sigma = 2.5$ Myr, corresponding to an FWHM exceeding 5 Myr. Considering the very young ages measured in this luminosity interval, and the fact that, for our modeling, the minimum age included in our simulation is 0.5 Myr, this finding implies that the best-fit age distribution for the luminous part of the PMS sequence is highly skewed, with a peak at $\sim 1$ Myr and a broad tail toward older ages. On the other hand, for the faintest end of our stellar sample, our method finds that a more modest age spread is needed to produce a simulated distribution of stars as broad as the observed sequence.

The differences in the derived age spreads as a function of luminosities may be due to an incorrect modeling of the two-dimensional isochrones we have used. In particular, some of the broadening effects we included in the model might be mass dependent. For example, variability and accretion might affect more the brightest members, and less the smallest masses, whereas we assumed identical distributions of these effects along the entire mass range. Considering our results, and averaging the age spreads at different luminosities, the average age spread for the entire stellar sample, derived in Section 6.2 should be a fairly well approximation of the real spread.

### Table 4

| Magnitude Range | Age (Myr) | Age Spread ($\sigma$) (Myr) | Age (Myr) | Age Spread ($\sigma$) (Myr) |
|-----------------|-----------|----------------------------|-----------|----------------------------|
| $22 \leq m_{555} \leq 23$ | 0.9 | 2.2 ± 0.7 | 1.0 | 2.8 ± 0.7 |
| $23 \leq m_{555} \leq 24$ | 1.4 | 1.7 ± 0.5 | 1.8 | 2.4 ± 0.9 |
| $24 \leq m_{555} \leq 25$ | 2.6 | 0.6 ± 0.6 | 2.7 | 0.9 ± 0.8 |
| $25 \leq m_{555} \leq 26$ | 3.6 | 0.8 ± 0.4 | 3.6 | 0.6 ± 0.4 |
| $26 \leq m_{555} \leq 27$ | 4.6 | 1.0 ± 0.2 | 4.7 | 0.0 ± 0.8 |
If we trust our models, the age–luminosity dependence we measure implies that stars in the low-mass regime have been formed some Myr before the intermediate-mass members. However, this is not what is observed in young Galactic 1–2 Myr old associations, where no evidence of highly biased IMF are found (Hartmann 2003). An explanation for this age difference among different stellar masses is the assumed binarity. An overestimated binary fraction can lead to overestimation of the age for the low-mass stars, because the large numbers of unresolved low-mass binaries produce systematically wider modeled distributions of low-mass stars in the synthetic CMDs. We have investigated this effect by computing the average age in the five selected magnitude intervals (as those in Table 4), but with two-dimensional modeled distributions constructed assuming no binaries ($f = 0$). While in this case the mass–age dependence reduces, it does not vanish, with a reduced variation in the derived ages of the order of 1 Myr. Naturally, this is an extreme case since such a very low binary fraction is not realistic, but unfortunately the true binary fraction in the considered mass ranges cannot be directly resolved at the distance of the LMC. \textsuperscript{2}

We have investigated more in detail if the mass–age correlation we have found can be a biased result to an inaccuracy in the modeling of the two-dimensional isochrones. In particular, since the effect caused by accretion veiling is the most uncertain among those we have considered (mostly because of the lack of mass accretion estimates for LH 95), we have tried to redo the entire analysis presented in this paper excluding accretion from our modeling. In case of no accretion, we find that the age of the cluster is older (of the order of 6 Myr, depending on the other parameters), the age spread is larger (several Myr), and still, there is consistent mass–age correlation, with the very low mass stars systematically older than the intermediate-mass stars.

Taking into account the above, and if we assume that indeed such a dependence of age with mass should not be observed, as in Galactic young clusters, we can attribute the observed dependence of age with mass to an inaccuracy of the PMS models we have utilized. Actually, the inconsistency between the ages predicted by the FRANEC and Siess isochrones suggests that the PMS stellar modeling is still far from completely reliable. For a metallicity close to solar, typical of the Galaxy, several studies have been presented to constrain the evolutionary models based on empirical isochrones (e.g., Mayne et al. 2007). Certainly new data that sample different young associations and clusters in the metal-poor environment of the MCs will be exceptionally useful to test the agreement of data with the PMS models also for lower metallicities.

7. DISCUSSION

We stress that our results on the age determination of LH 95 and the evidence for an age spread in the system are valid under the assumption of the correctness of the sources of spread in the CMD considered in our models (Section 5.1). It should be noted that having included, apart from photometric uncertainties, several important biasing factors, namely, differential reddening, stellar variability, accretion binarity, and crowding, our modeling technique is more accurate than similar previous studies, which completely neglect these biases. Using our models, we cannot reproduce the observed sequence assuming a coeval population, but an additional spread in age is measured. Nonetheless, as we have shown, our results are very sensitive on the assumptions used to model observed populations in the CMD. In particular, changes in the assumed binary fraction alter significantly both the average cluster age and the derived age spread. Also, as discussed previously, the overall effect of accretion veiling influences both this quantities. For our modeling, we assumed accretion and variability measurements from the ONC. If ongoing accretion in LH 95 is weaker than we assumed, for instance because of an earlier disruption of the circumstellar disks caused by the flux of the early-type stars present in the region (three of which are O-type stars; N. Da Rio et al. 2010, in preparation) or simply because of the older age, the measured age would turn out be older and the age spread larger that the results we have presented.

Slow star formation scenarios have been proposed based on both observational and theoretical evidence (Palla & Stahler 2000; Palla et al. 2005; Tan et al. 2006); the loose, not centrally concentrated geometry of LH 95, and its relatively low stellar density and mass (see Paper I) would be compatible with a slow conversion of gas into stars (Kennicutt 1998). Unfortunately, the absence of dynamical information on our region, as well as a precise estimate of the gas mass do not allow us to infer much about the past evolution of LH 95.

On the other hand, additional unconsidered effects that can displace PMS stars in the CMD may still be present and affect the derived age spread; we briefly discuss them here. The presence of circumstellar material can cause star light to be scattered in the line of sight direction (e.g., Kraus & Hillenbrand 2009; Guarcello et al. 2010). This effect may either increase the observed optical flux with respect to that of the central star, or reduce it, if the central star is partially obscured by the circumstellar material. This scenario, however, requires both the presence of a significant circumstellar environment and a particular geometrical orientation, so, if present, would affect only a fraction of the stars in a cluster. In any case, as shown in Guarcello et al. (2010), stars observed in scattered light appear significantly bluer than the PMS sequence. If such sources are present in our photometric catalog, they would affect only the old tail of our age distribution, or even be excluded as candidate MS field stars. Thus, we are confident that neglecting this effect should not affect significantly our findings. Another unconsidered effect that could bias the measured age spread is that presented by Baraffe et al. (2009). Under certain conditions, if a star accretes a significant amount of material it cannot adjust its structure quickly enough, so appears underluminous for a star of its mass. This effect will have a memory, which may persist for several Myr according to the authors, and will cause some young stars to appear older than they are. This would also cause an additional apparent spread in the H–R diagram even in case of a coeval population. The scenario proposed by Baraffe et al. (2009), however, requires episodes of vigorous mass accretion, with a high $M$. With no direct measurements of accretion in LH 95, and no information about the past history of accretion, as a function of mass, we are unable to estimate if this scenario could play any role in our region, and if so, to what extent.

8. SUMMARY

In this study, we utilize the photometric catalog of sources in the young LMC stellar association LH 95, based on HST/ACS photometry in F555W and F814W filters (corresponding to standard $V$ and $I$ bands). The catalog, which covers stellar masses as low as 0.2 $M_{\odot}$, was used in Paper I to derive for the first time the subsolar IMF in the LMC. Here, we focus on the age of the system and the determination of the age spread from the CMD broadening of its PMS stars. We first refine the

\textsuperscript{2} At the distance of the LMC, 1 ACS/WFC pixel corresponds to $\sim$2400 AU.
decontamination of our sample from field stars by eliminating any remaining field MS stars still present in the first catalog of true stellar members (Paper I). Our photometric catalog of member stars of LH 95 covers ~2000 members in the PMS and UMS regions of the CMD.

We develop a maximum-likelihood method, similar to that described in Naylor & Jeffries (2006) to derive the age of the system accounting simultaneously for photometric errors, unresolved binarity, differential extinction, accretion, stellar variability, and crowding. The application of the method is performed in several steps: first, a set of PMS isochrones is converted to two-dimensional probability distributions in the magnitude–magnitute plane; one model distribution is constructed for each considered age. A well-populated \((N > 10^6 \text{ stars})\) synthetic population is thus created by applying the aforementioned sources of displacement in the CMD—except for photometric errors—to the isochrone for each chosen age. A maximum-likelihood method is then applied to derive the probability, for each observed star to have a certain age, taking its photometric uncertainty into account. The multiplication of the probability, for each observed star to have a certain age, taking its photometric uncertainty into account. The multiplication of the probability, for each observed star to have a certain age, taking its photometric uncertainty into account.

We apply our method to the stellar catalog of LH 95, assuming the IMF and reddening distribution for the system, as they were derived in Paper I. We assume distributions of the variations in both \(m_{555}\) and \(m_{814}\) magnitudes caused by accretion and variability due to dark spots by utilizing data for the ONC from the literature. We model age-dependent probability distributions using two sets of evolutionary models, that by Siess et al. (2000), for the metallicity of the LMC, and a new grid computed with the FRANEC stellar evolution code (Chieffi & Straniero 1989; Degl’Innocenti et al. 2008) for the same metallicity. We evaluate the accuracy of our age estimation and we find that the method is fairly accurate in the PMS regime, while the precision of the measurement of the age is lower at higher luminosities and old ages. Our method cannot constrain the age of UMS stars due to the small dependence of CMD position with age. Our treatment showed that the age determination is very sensitive to the assumed evolutionary model and binary fraction; the age of LH 95 is found to vary from \(\sim 2.8 \text{ Myr} \sim 4.4 \text{ Myr}\), depending on these factors. Assuming \(f = 0.5\), the average binary fraction in our mass range (Lada 2006), the best-fit age of the system derived for this binary fraction is determined to be of the order of 4 Myr. Despite the sparse and clumpy spatial distribution of PMS stars in LH 95, we find no evidence of systematic variations of age with location.

With our method we are also able to assess the true appearance of an age spread among the PMS stars of LH 95, as an indication that star formation in the system occurred not spontaneously but continuously. Our treatment allowed us to disentangle any real age spread from the apparent CMD spread caused by the physical and observational biases which dislocate the PMS stars from their theoretical positions in the CMD (see Section 5.1). We find that LH 95 hosts an overall age spread well modeled by a Gaussian distribution with \(FWHM = 2.8–4 \text{ Myr}\).

We measure separately ages and age spread as a function of luminosity. For the intermediate masses \((M \sim 1 M_\odot)\), we find younger average ages and larger age spread, whereas for low-mass stars \((M \sim 0.1 M_\odot)\), older ages and more modest age spreads. We assign this dependence of age with mass not to a real effect but to imperfections of the PMS evolutionary models, which tend to predict lower ages for the intermediate masses, and higher ages for low-mass modeling. In any case, under the validity of the assumptions used in our modeling, we rule out the possibility that the population of LH 95 is coeval.

Finally, we remark that, though the method we have presented provides an exceptional tool for an unbiased investigation of ages and age spreads in young PMS regions, the results are highly dependent on the modeling of the assumed sources of apparent spread in the CMD. Only a precise characterization of each of these effects, in different regions, will enable us to give the final answer on the long-standing questions about the duration of star formation.

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