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Specific surface area of polydispersions as a function of size distribution sharpness

PAULO FILIPE T. LOPES, JOSÉ AURÉLIO M. DA LUZ & FELIPE O. MILHOMEM

Abstract: Knowledge of particulate system properties is very important in various industrial instances and the possibility of fast predicting the behavior of such systems is an important control tool. The specific surface area of simulated particulate systems was studied as a function of the sharpness parameter of the size distribution of the Gates–Gaudin–Schumann, Gaudin–Meloy and Rosin–Rammler equations. The results showed good statistical adherence, especially the Rosin–Rammler equation, in situations where it is the best descriptor of particle size distribution.

Key words: particle size distribution, sharpness index, statistical distribution, surface area.

INTRODUCTION

Specific surface area of granular media (bulk material in general) plays a relevant role in various instances of engineering, such as reactors with solid/liquid, liquid/gas, or solid/gas systems. However, this parameter is generally difficult to quantify, especially when speed in its determination is required (as in the case of its use as input in algorithms for dynamic online process control). Instrumental techniques, such as N₂ adsorption (in accordance with the BET isotherm), liquid porosimetry, or X-ray tomography, are usually costly, time-consuming and sometimes unfeasible in the environment of a production process. Therefore, indirect and quick techniques for determining the specific surface area become very attractive. The development of a tool for inference of this parameter, based on the particle size distribution of the particulate medium, was the motivation for the present work.

Particle size distribution of granular media can commonly be described appropriately by equations, with behavior similar to a lognormal particle size distribution (Gaussian curve for argument logarithm) (Allen 1981, Bender 2000). Generally, however, other statistical distributions are used, either for greater statistical adherence or for more convenient algebraic manipulation simplicity (since distributions derived from the Gaussian distribution are not Riemann integrable).

Statistical distributions currently used in the analysis of particulate systems (in various instances such as minerometallic, chemical, ceramic, food and other processes) are, for example, Gates–Gaudin–Schumann (GGS), Gaudin–Meloy (GM) and Rosin–Rammler (RR). They are given by equations 1, 2 and 3 of Table I.

In addition to the knowledge on its own of the particle size characteristics and its influence on size-dependent operations, the data obtained from such features can be used in other instances, such as relating the friction angle to the bulk material’s size distribution. (Edil et al. 2007), predicting soil water retention (Mohammadi & Meskini-Vishkaee 2012, Nasta et
Table I. Usual distributions for characterization of polydisperse systems (real argument \( x \) is particle size).

| Tag | Accumulated passing \( (Y(x)) \) | Parameters | Eq. |
|-----|----------------------------------|------------|-----|
| GGS | \( Y(x) = \left( \frac{x}{x_{\text{max}}} \right)^n \) | \( n \) — distribution exponent; \( x_{\text{max}} \) — top size. | (1) |
| GM  | \( Y(x) = 1 - \left( 1 - \frac{x}{x_{\text{max}}} \right)^n \) | \( n \) — distribution exponent; \( x_{\text{max}} \) — top size. | (2) |
| RR  | \( Y(x) = 1 - \exp \left[ \frac{1}{2} \times \left( \frac{x}{x_{50}} \right)^n \right] \) | \( n \) — sharpness (uniformity coefficient); \( x_{50} \) — median diameter. | (3) |

Another interesting example is the interrelationship between specific surface of aggregates and their degree of packing, with impact on concrete properties such as mortar workability and strength of cured concrete (Hunger & Brouwers 2009). The specific surface area computed from the measured particle size distributions and its corresponding value estimated by Blaine permeability has revealed a constant ratio and could be used to predict the quality of the concrete. Ultimately, the specific surface area and particle shape affect the flowability and water requirement of Portland cement mortar. Working on this line, Ghasemi et al. (2018) calculated the specific area of granular systems as a function of particle size curves and particle morphometric modeling (assuming polyhedral particle shapes).

In the present study, the equations in Table I were used to simulate particle size distributions of hypothetical particulate systems by varying the sharpness parameter \( (n) \) and correspondingly calculating the resulting volumetric specific surface area \( (s_v) \). Thus, it was possible to generate a cloud of ordered \( n \) versus \( s_v \) pairs, of sufficient magnitude to allow the mathematical prediction of volumetric specific surface area of particulate systems only on the basis of statistical parameters of their distributions descriptors, as discussed below.

**MODELING THE PROBLEM**

The specific volumetric surface area of a particle is the ratio between its area and its volume. For spherical particles the solution is trivial. In the case of morphologically irregular particles, one can consider the equivalence of the particle to a sphere of the same dimensions (in terms of diameter, projection area, volume, etc.), by adopting appropriate morphometric coefficients (for sphere: \( k_s = k_v = 1 \)). The specific volume surface area \( s_v \) of an irregular particle (from a size class \( i \)) is then equal to:

\[
s_v = \frac{\text{particle surface area}}{\text{particle volume}} = A \cdot \frac{x_i^2}{x_i^3} \cdot k_s \cdot \frac{\pi x_i^2}{6} \cdot k_v \cdot \frac{6}{x_i} = \frac{A \cdot \pi x_i^2}{6} \cdot k_s \cdot k_v \cdot \frac{6}{x_i} \quad (4)
\]
This demonstrates the dependence of volumetric surface area on particle size ($x$), tending to very low values for very large particles and infinite value for very tiny particles.

It is convenient to mention morphometric coefficients often vary with the particle size class. A classic compilation by Orr Jr. & DallaValle (1959) can be cited to illustrate this phenomenon, shown in Table II and Figure 1, where the values of morphometric coefficients of quartz sand and crushed quartz are compared. Quartz particles (mineral free of cleavage planes) subjected to natural transport and physical weathering processes tend to decrease their sphericity with increasing size. On the other hand, artificially comminuted particles increase their irregularities with increasing degree of comminution (smaller particles tend to be more irregular).

The masses of particles retained in each size class, $i$ (where $x_i > x_{i+1}$), can be calculated from the frequency ($f_i$), whereas the approximate number of particles ($N_i$) is obtained by dividing the retained volume in each class by the volume of the average particle of this class. Thus, it results:

$$f_i = Y(x_i) - Y(x_{i+1}) = \frac{m_i}{m_i} \cdot m_i = m_i \times f_i$$

$$N_i = \frac{\frac{m_i}{\rho}}{\frac{k_i \times \frac{\pi}{6} \times x_i^3}{}}$$

$$N_i = \frac{6 \times \frac{m_i}{k_i \times \frac{\pi}{6} \times x_i^3}}{m_i \times f_i} \cdot N_i \times k_i = 6 \frac{\pi \times f_i}{x_i^3}$$

Table II. Surface and volume shape coefficient for some particulate materials (modified data from Orr Jr. and DallaValle’s).

| Granular system | $x_50$ [m] | Shape coefficient | Primary data source (apud Orr Jr. & DallaValle, 1959) |
|-----------------|------------|-------------------|-----------------------------------------------|
| White sand      | 2.90 x 10^{-3} | $k_s = 0.668$    | DallaValle (1938)                            |
| White sand      | 8.00 x 10^{-4} | $k_s = 0.859$    | DallaValle (1938)                            |
| White sand      | 4.00 x 10^{-4} | $k_s = 0.828$    | DallaValle (1938)                            |
| Filter sand     | 6.00 x 10^{-4} | $k_s = 0.859$    | DallaValle (1938)                            |
| Filter sand     | 5.00 x 10^{-4} | $k_s = 0.923$    | DallaValle (1938)                            |
| Crushed quartz  | 8.30 x 10^{-3} | $k_v = 0.796$    | Martin (1927–8)                              |
| Crushed quartz  | 4.90 x 10^{-3} | $k_v = 0.796$    | Martin (1927–8)                              |
| Crushed quartz  | 1.10 x 10^{-3} | $k_v = 0.700$    | Martin (1927–8)                              |
| Crushed quartz  | 3.00 x 10^{-4} | $k_v = 0.668$    | Martin (1927–8)                              |
| Crushed quartz  | 2.20 x 10^{-3} | $k_v = 0.535$    | Martin (1927–8)                              |
| Crushed quartz  | 1.90 x 10^{-3} | $k_v = 0.325$    | Goldman & DallaValle (1939)                   |
| Crushed quartz  | 8.90 x 10^{-4} | $k_v = 0.286$    | Goldman & DallaValle (1939)                   |
| Crushed quartz  | 4.90 x 10^{-4} | $k_v = 0.516$    | Martin (1927–8)                              |
| Crushed quartz  | –           | $k_v = 0.267$    | Hatch & Choate (1929)                         |
| Calcite         | –           | $k_v = 0.258$    | Hatch & Choate (1929)                         |
| Hornblende      | 5.00 x 10^{-3} | $k_v = 0.038$    | Goldman & DallaValle (1939)                   |
| Dolomite        | 4.00 x 10^{-3} | $k_v = 0.363$    | Goldman & DallaValle (1939)                   |
| Feldspar        | 4.00 x 10^{-3} | $k_v = 0.497$    | Goldman & DallaValle (1939)                   |
Substituting equation (10) into (4) finally results:

$$s_i = 6 \times \sum_{i=1}^{m} k_i \times f_i = s_m \times \rho$$

In the previous equation $s_m$ is the specific mass surface area (expressed in m²/kg) and $\rho$ is the particle’s density (expressed in kg/m³). The simulations were performed based on the above, adopting spherical particles ($k_v = k_s = 1$), and generating a size distribution with relative size ranging from 0 to 1 (with increments of 0.0025).

The mass cumulative pass-through proportion ($Y(x)$) was simulated for distributions described by GGS, GM and RR equations. The exponent values of the distributions were in the range: $0.001 \leq n \leq 100$. Although such extreme limits are not usual in industrial reality, they were taken into account for a broader range of results. Based on cumulative mass function, the retained fraction in each size class was determined and, accordingly, the specific surface area of each fraction was calculated. The specific surface areas for the simulated distributions were obtained from the sum corresponding to all granulometric classes.

Once in possession of the specific surface area values for the tested values of $n$, the regression between sharpness parameter and volumetric surface area was performed, thus making it possible to obtain a general equation to predict the surface area based on $n$.

**RESULTS**

Figure 2 illustrates the surface area curves versus the exponent of the distribution. Similar behavior of Gates–Gaudin–Schumann and Rosin–Rammler distributions was observed: small initial values, for small values of $n$, which increase until they decrease again, when the exponent is greater than 1, decreasing asymptotically. In contrast, for Gaudin–Meloy distribution a gradual increase in surface area is observed as the exponent increases, tending to asymptotic saturation.
A regression analysis was performed from the simulated values, using nonlinear regression with Easyplot software. Table III systematizes the predicted equations of the specific surface area of polydisperse spheroidal particulate systems as a function of the distribution exponent. Using literature results (da Luz 2005), one can immediately calculate the surface area of the granular system which is well described by the modified logistic equation (de Lynch-Rao's equation), shown below (Eq. 13).

$$\beta = 34.555 \times e^{\left( \frac{-3.9}{n^{0.5984}} \right)}$$

Another aspect to be highlighted is that in Gaudin–Meloy and Gates–Gaudin–Schumann distributions there is a confounding effect between the central tendency parameter and the dispersion parameter, both embedded in the distribution exponent (since the second parameter in the equation refers to the top size). Thus, concomitant change in the distribution median occurs when parameter \( n \) is modified, logically indicating the exponent cannot be interpreted as an isolated measure of the degree of uniformity of sizes (often referred to as sharpness). In contrast, Rosin–Rammler distribution shows total independence of the two parameters, since they are unambiguously explicit in the equation, which facilitates the analysis of the behavior of particulate systems occurring in industrial practice. By the way, the usual range of sharpness in industry applications is: \( 0.5 \leq n \leq 4.0 \).
Table III. Equations and parameters from nonlinear regression.

| Distribution | Regression equations | Fitting parameters | Equation |
|--------------|----------------------|--------------------|----------|
| GM           | \( s_v = a \times \left[1 - \exp\left(-b \times n^c\right)\right] \) | \( a = 1022.79 \) \( b = 0.0385428 \) \( c = 0.808936 \) | (10) |
| GGS          | \( s_v = \left[ a \times k \times b^a \times n^{(a+b)^2} \right] + c \) | \( a = 1.9592 \) \( b = 0.38077 \) \( c = 6.4148 \) \( k = 144.129 \) | (11) |
| RR           | \( s_v = a(n) \times d_{50}^{-b(n)} \) | \( a(n) = \frac{290.056 \times n^{1.65813}}{n^{1.65813} + 0.20708^{1.65813}} + 6 \) | (12a) |
|              |                      | \( b(n) = 1 - \exp\left(-\left(\frac{n}{0.67235}\right)^{1.6902}\right) \) | (12b) |

CONCLUSIONS

The equations obtained here for the prediction of the specific surface area of particulate systems accurately described by the Rosin–Rammpler, Gates–Gaudin–Schumann and Gaudin–Meloy statistical distributions (and, furthermore, by modified logistic distribution) allow the rapid implementation of process monitoring and control algorithms, in various instances of engineering, where specific surface area plays a relevant role, such as chemical reactors with solid/liquid, liquid/gas, or solid/gas systems, filtration circuits, solvent extraction processing routes, formulation of Portland cement mortars, adsorption processes, and fines agglomeration operations, just to name a few.

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