Hot Corinos Chemical Diversity: Myth or Reality?

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Abstract

After almost 20 years of hunting, only about a dozen hot corinos, hot regions enriched in interstellar organic molecules (iCOMs), are known. Of them, many are binary systems with the two components showing drastically different molecular spectra. Two obvious questions arise. Why are hot corinos so difficult to find and why do their binary components seem chemically different? The answer to both questions could be a high dust opacity that would hide the molecular lines. To test this hypothesis, we observed methanol lines at centimeter wavelengths, where dust opacity is negligible, using the Very Large Array interferometer. We targeted the NGC 1333 IRAS 4A binary system, for which one of the two components, 4A1, has a spectrum deprived of iCOMs lines when observed at millimeter wavelengths, while the other component, 4A2, is very rich in iCOMs. We found that centimeter methanol lines are similarly bright toward 4A1 and 4A2. Their non-LTE analysis indicates gas density and temperature (≥2 × 10^6 cm^-3 and 100–190 K), methanol column density (∼10^19 cm^-2), and extent (∼35 au in radius) similar in 4A1 and 4A2, proving that both are hot corinos. Furthermore, the comparison with previous methanol line millimeter observations allows us to estimate the optical depth of the dust in front of 4A1 and 4A2, respectively. The obtained values explain the absence of iCOMs line emission toward 4A1 at millimeter wavelengths and indicate that the abundances toward 4A2 are underestimated by ∼30%. Therefore, centimeter observations are crucial for the correct study of hot corinos, their census, and their molecular abundances.

Unified Astronomy Thesaurus concepts: Interstellar medium (847); Protostars (1302); Star formation (1569); Chemical abundances (224); Astrochemistry (75)

1. Introduction

Interstellar complex organic molecules (iCOMs) are molecules detected in the interstellar medium containing carbon and at least six atoms (Herbst & van Dishoeck 2009; Ceccarelli et al. 2017). These molecules are of particular interest because they carry a substantial fraction of carbon that can be used for prebiotic chemistry (e.g., Caselli & Ceccarelli 2012).

In solar-like young Class 0 protostars, iCOMs are found in relatively large quantities toward the so-called hot corinos, which are compact (<100 au), hot (>100 K), and dense (≥10^7 cm^-3) regions enriched in iCOMs at the center of the envelopes accreting the future star (Ceccarelli 2004; Ceccarelli et al. 2007; Caselli & Ceccarelli 2012).

The first hot corino was discovered in 2003 toward the Class 0 source IRAS 16293–2422 (e.g., Cazaux et al. 2003; Jørgensen et al. 2016; Manigand et al. 2020). Since then other Class 0 hot corinos have been discovered: NGC 1333 IRAS 4A (hereafter IRAS 4A; e.g., Bottinelli et al. 2004; Taquet et al. 2015; De Simone et al. 2017; López-Sepulcre et al. 2017; Sahu et al. 2019), NGC 1333 IRAS 2A, NGC 1333 IRAS 4B (e.g., Jørgensen et al. 2005; Bottinelli et al. 2007; Maury et al. 2014; De Simone et al. 2017), HH 212 (Codella et al. 2016; Bianchi et al. 2017; Lee et al. 2017, 2019), B335 (Imai et al. 2016), L483 (Oya et al. 2017; Jacobsen et al. 2019), Barnard1b-S (Marcellino et al. 2018), Ser-emb 1 (Martin-Domenech et al. 2019), and BHR71-IRSI (Yang et al. 2020). Lately, a few more evolved Class 1 hot corinos were also discovered: NGC 1333 SVS13A (De Simone et al. 2017; Bianchi et al. 2019), B1a (Öberg et al. 2014), and Ser-emb 17 (Bergner et al. 2019). Therefore, after almost 20 years, only about a dozen hot corinos are known. Recent surveys concluded that ∼30% of low-mass Class 0/I protostars show emission from at least three iCOMs (De Simone et al. 2017; Belloche et al. 2020).

Most of the hot corinos cited above turn out to be binary systems when imaged at high angular resolution. This is in agreement with previous surveys that found that 40–60% of protostars are multiple systems (Maury et al. 2014; Tobin et al. 2016). Interestingly, with the first hot corino maps it became
clear that the two objects in a given binary system can substantially differ in molecular complexity. Illustrative examples are provided by IRAS 16293–2422 and IRAS 4A (Jørgensen et al. 2016; López-Sepulcre et al. 2017). IRAS 16293–2422 is composed of two sources, A and B, separated by 5′1 (~720 au), where source A, weaker in millimeter continuum emission, is brighter in iCOMs lines than source B (e.g., Caux et al. 2011; Pineda et al. 2012; Jørgensen et al. 2016; Manigand et al. 2020). IRAS 4A, located in the NGC 1333 region in the Perseus cloud at a distance of (299 ± 15) pc (Zucker et al. 2018), is also a binary system composed of IRAS 4A1 and IRAS 4A2 (hereafter 4A1 and 4A2), separated by 1′8 (~540 au): while 4A1 is brighter in the millimeter continuum than 4A2, only 4A2 shows bright iCOMs lines (Taquet et al. 2015; De Simone et al. 2017; López-Sepulcre et al. 2017). However, not always the brightest millimeter continuum source in a binary system is the one weak iCOMs emission (see, e.g., Ospina-Zamudio et al. 2018).

In summary, despite two decades of hunting, only a dozen hot corinos are known so far. Of them, many are binary systems with the two components showing drastically different molecular spectra. Two related questions arise: (1) Why are hot corinos so difficult to find? While it is known that not all Class 0/I sources possess hot corinos (e.g., Sakai & Yamamoto 2013; Higuchi et al. 2018; Bouvier et al. 2020), observational biases might hamper their detection. (2) Why do coeval objects seem drastically different in their chemical composition? Is this a real difference or is it only/mostly due to observational biases?

A major observational bias could be caused by the dust opacity, which could be very high in Class 0/I sources, due to their high densities and, consequently, column densities (e.g., Miotello et al. 2014; Galván-Madrid et al. 2018; Galametz et al. 2019). If the effect of dust absorption is not negligible, there are three major consequences: (1) hot corinos may be difficult to detect in the millimeter (also) because of the high dust absorption of the iCOMs lines; (2) the molecular complexity diversity observed in binary system objects may reflect a difference in the front dust column density rather than a real chemical difference of the two objects; (3) the iCOM abundances in hot corinos could have been underestimated so far. In order to test this hypothesis, we targeted the IRAS 4A binary system, where the two objects show extremely different iCOM line spectra at millimeter wavelengths (see above), and carried out observations of several methanol lines, one of the simplest iCOMs, at centimeter wavelengths where the dust is optically thin.

2. Observations

The IRAS 4A system was observed at 1.3 cm using K-band receivers (18–26.5 GHz) of the Karl G. Jansky Very Large Array (VLA) in C-configuration (35–3400 m) on 2018 December 10 (project ID: VLA/18B-166). We targeted 10 CH3OH lines, with frequencies from 24.9 to 26.4 GHz, upper level energies $E_{up}$ from 36 to 175 K and Einstein coefficients $A_{fi}$ in the range (0.5–1.1) $\times 10^{-7}$ s$^{-1}$ (Table 1). The observed spectra were divided into eight spectral windows with $\sim 0.017$ MHz (0.2 km s$^{-1}$) spectral resolution and $\sim 1''$ (~300 au at the distance of IRAS 4A) angular resolution. The observations were centered on 4A2, at $\alpha(J2000) = 03^h29^m10.4^s, \delta(J2000) = 31^\circ13'32''$. The flux calibrators were J0137+3309 and J0521+1638, while the bandpass and the gain ones were J0319+4130 and J0336+3218, respectively. The absolute flux calibration error is $\leq 15\%$.14

The data reduction and cleaning process were performed using the CASA15 package while data analysis and images were performed using the GILDAS16 package. We obtained a continuum image by averaging line-free channels from all the spectral windows (Figure 1). We self-calibrated, in phase amplitude, using the line-free continuum channels and applied the solutions to both the continuum and molecular lines. The dynamic range, as defined by peak source flux over rms noise, was improved by 20% by the self-calibration. The final rms noise in the continuum image, $3 \mu$Jy beam$^{-1}$, is consistent with that reported by the VLA Exposure Time Calculator for a line-free continuum bandwidth of 4.5 GHz, 26 antennas, and an on-source integration time of 3 hr. The cube was subsequently continuum subtracted, smoothed to 1 km s$^{-1}$ (~0.08 MHz), and cleaned in CASA using a multiscale deconvolution17 (scales = [0, 5, 15, 18, 25]) with natural weighting. The synthesized beams for each spectral window are reported in Table 1. The half power primary beam is $\sim 80''$.

3. Results

3.1. Continuum Emission Map

Figure 1 reports the map of the continuum emission at 25 GHz. The two continuum peaks mark the two protostars, whose coordinates ($\alpha(J2000) = 03^h29^m10.536^s, \delta(J2000) = 31^\circ13'31.07''$, for 4A1, and $\alpha(J2000) = 03^h29^m10.43^s, \delta(J2000) = 31^\circ13'32.1''$, for 4A2) are consistent with those derived by Tobin et al. (2016) and López-Sepulcre et al. (2017) with higher angular resolution observations. Since the angular resolution of our observations ($\sim 3''$) is smaller than the separation between 4A1 and 4A2 ($\sim 8''$), they are clearly disentangled in our images, even if individually unresolved with the current resolution.

At centimeter wavelengths, 4A1 shows a brighter continuum emission (due to dust or free–free) than 4A2. The peak fluxes are $(2.1 \pm 0.3)$ mJy beam$^{-1}$ and $(0.47 \pm 0.07)$ mJy beam$^{-1}$ toward 4A1 and 4A2, respectively. Taking into account the slightly different wavelength (1.05 cm) and angular resolution ($\sim 0.1''$), these values are consistent with the ones measured by Tobin et al. (2016): $(1.3 \pm 0.2)$ mJy beam$^{-1}$ for 4A1 and $(0.38 \pm 0.04)$ mJy beam$^{-1}$ for 4A2.

3.2. Methanol Lines

All the targeted methanol lines are detected with a signal-to-noise ratio larger than 3 (Table 1). Their velocity-integrated spatial distribution is shown in Figure 1. The methanol emission peaks exactly toward the 4A1 and 4A2 continuum peaks, and it is well disentangled, even if unresolved at the current angular resolution, around the two protostars.

Figure 2 shows the 10 methanol line spectra, isolated and not contaminated by other species, extracted toward the 4A1 and 4A2 protostars.
Table 1
Spectral Parameters, Synthesized Beams, and Gaussian Fit Results of the CH$_3$OH Lines Extracted toward the 4A1 and 4A2 Continuum Peaks

| Transition   | Frequency$^a$ (GHz) | $E_{up}$ $^a$ (K) | log $A_0$ $^a$ | Synthesized Beam $^{maj \times min}$ (PA) ($'' \times ''$) | IRAS 4A1 | IRAS 4A2 |
|--------------|---------------------|-------------------|----------------|---------------------------------|---------|---------|
|              |                     |                   |                | $\int T_d dv$ $^b$ (K km s$^{-1}$) | $V_{peak}$ $^b$ (km s$^{-1}$) | FWHM$^b$ (km s$^{-1}$) | rms$^c$ (K) | $\int T_d dv$ $^b$ (K km s$^{-1}$) | $V_{peak}$ $^b$ (km s$^{-1}$) | FWHM$^b$ (km s$^{-1}$) | rms$^c$ (K) |
| 3(2, 1)–3(1, 2) E | 24.92871            | 36                | −7.2           | 0.97 $\times$ 0.95 (−12)         | 17(4)   | 6.5(0.2) | 2.7(1.2) | 1.0     | 34(3)   | 6.8(0.2) | 3.1(0.3) | 0.9     |
| 4(2, 2)–4(1, 3) E | 24.93347            | 45                | −7.1           | 0.97 $\times$ 0.95 (−12)         | 23(3)   | 6.5(0.2) | 3.7(0.5) | 1.0     | 32(3)   | 6.9(0.2) | 3.1(0.3) | 0.9     |
| 2(2, 0)–2(1, 1) E | 24.93438            | 29                | −7.2           | 0.97 $\times$ 0.95 (−12)         | 19(3)   | 6.5(0.2) | 3.0(0.6) | 1.0     | 27(3)   | 6.9(0.2) | 2.9(0.3) | 0.9     |
| 5(2, 3)–5(1, 4) E | 24.95908            | 57                | −7.1           | 0.97 $\times$ 0.95 (−12)         | 19(3)   | 6.1(0.3) | 4.2(0.9) | 0.9     | 32(3)   | 6.8(0.2) | 3.1(0.3) | 0.9     |
| 6(2, 4)–6(1, 5) E | 25.01812            | 71                | −7.1           | 0.97 $\times$ 0.95 (−19)         | 19(3)   | 6.3(0.3) | 3.5(0.6) | 1.0     | 31(3)   | 6.7(0.1) | 2.5(0.2) | 1.0     |
| 7(2, 5)–7(1, 6) E | 25.12487            | 87                | −7.1           | 0.98 $\times$ 0.95 (−21)         | 20(2)   | 7.0(0.2) | 3.2(0.5) | 0.9     | 35(3)   | 6.8(0.2) | 2.9(0.3) | 1.2     |
| 8(2, 6)–8(1, 7) E | 25.29442            | 106               | −7.0           | 0.96 $\times$ 0.94 (−11)         | 13(3)   | 6.8(0.3) | 3.4(0.9) | 0.9     | 32(2)   | 6.8(0.1) | 2.8(0.3) | 0.9     |
| 9(2, 7)–9(1, 8) E | 25.54140            | 127               | −7.0           | 0.96 $\times$ 0.92 (−50)         | 18(2)   | 6.5(0.2) | 2.7(0.5) | 0.8     | 31(2)   | 6.9(0.1) | 2.6(0.2) | 0.9     |
| 10(2, 8)–10(1, 9) E | 25.87827             | 150              | −7.0           | 0.97 $\times$ 0.93 (−35)         | 20(2)   | 6.4(0.2) | 3.2(0.5) | 0.8     | 32(2)   | 6.9(0.1) | 2.6(0.2) | 0.8     |
| 11(2, 9)–11(1, 10) E | 26.31312            | 175               | −6.9           | 0.94 $\times$ 0.91 (−35)         | 24(4)   | 6.0(0.4) | 4.9(0.9) | 1.1     | 31(3)   | 6.8(0.1) | 2.9(0.3) | 0.9     |

Notes.
$^a$ Spectroscopic parameters by Xu et al. (2008) are from the Cologne Database for Molecular Spectroscopy (Müller et al. 2005).
$^b$ Results of the Gaussian fit algorithm.
$^c$ The rms is computed over each spectral window.
The lines are slightly brighter toward 4A2 than 4A1, whereas the linewidths are very similar (see also Table 1). We derived the velocity-integrated line intensities for each detected CH$_3$OH transition using a Gaussian fit, being the profile Gaussian-like. The fit results for both sources, namely the integrated emission ($\int I_v dv$), the linewidth (FWHM), the peak velocities ($V_{\text{peak}}$), and the rms computed for each spectral window, are reported in Table 1. The velocity peaks are consistent with the systemic velocity of the molecular envelope surrounding IRAS 4A ($\sim 6.7$ km s$^{-1}$; Choi 2001). The line-widths are between 3 and 4 km s$^{-1}$ in agreement with those found by Taquet et al. (2015) and López-Sepulcre et al. (2017) toward 4A2 at millimeter wavelength.

In summary, our new VLA observations show a first clear important result: the detection of methanol emission toward 4A1, the protostar where previous millimeter observations showed no iCOMs emission (López-Sepulcre et al. 2017).

4. Centimeter versus Millimeter Observations: Dust Absorption Derivation

We compared our new centimeter observations of methanol lines with previous ones at 143–146 GHz in order to understand whether the dust absorption, more important at millimeter than at centimeter wavelengths, may explain the absence of iCOM millimeter line emission in 4A1 (López-Sepulcre et al. 2017). We first carried out a non-LTE analysis of the centimeter methanol lines from which we derived the gas temperature, density and CH$_3$OH column density toward 4A1 and 4A2 (Section 4.1). Then, using the same parameters, we predicted the methanol line intensities at 143–146 GHz, the frequency of the observations by Taquet et al. (2015; Section 4.2). Finally, we compared the predicted and measured millimeter line intensities and we attributed the difference to the absorption of the dust between us and the gas emitting methanol, via the usual equation:

$$I_{\nu}^{\text{obs}} = I_{\nu}^{\text{pred}} e^{-\tau_{\nu}}$$

in order to derive the dust optical depth toward 4A1 and 4A2, respectively (Section 4.3). Please note that the foreground dust opacity obtained by Equation (1) assumes that the absorbing dust fully covers the emitting gas area, which may not necessarily be the case. Yet, the derived attenuation of the methanol line intensities is still valid, even though it is only an average over the emitting gas area.

4.1. Non-LTE Analysis of the Centimeter Methanol Lines

To derive the physical properties of the gas emitting CH$_3$OH, namely gas temperature, density and methanol column density, we performed a non-LTE analysis using a Large Velocity Gradient (LVG) code (Ceccarelli et al. 2003). CH$_3$OH can be identified in A- and E-type due to the total spin ($I$) state of the hydrogen nuclei in the CH$_3$ group: A-type if the total spin function is symmetric ($I = 3/2$), E-type if...
asymmetric \( I = 1/2 \) \cite{Rabli2010}. We used the collisional coefficients of both types of CH$_3$OH with para-H$_2$, computed by Rabli & Flower \citeyear{Rabli2010} between 10 and 200 K for the first 256 levels and provided by the BASECOL database (Dubernet et al. 2013). We assumed a spherical geometry to compute the line escape probability (de Jong et al. 1980), the CH$_3$OH-A/CH$_3$OH-E ratio equal to 1, the H$_2$ ortho-to-para ratio equal to 3, and that the levels are populated by collisions and not by the absorption of the dust background photons whose contribution is very likely negligible due to the low values of the CH$_3$OH Einstein coefficients \( A_i \). Please note that the present LVG analysis only accounts for the line optical depth (to have also the dust \( \tau \) in the methanol emitting region would require information on the structure of the region that we do not have, as the emission is unresolved).

We ran a large grid of models (\( \geq 10^6 \)) covering the frequency of the observed lines, a total (CH$_3$OH-A plus CH$_3$OH-E) column density \( N_{\text{CH$_3$OH}} \) from \( 2 \times 10^{10} \) to \( 8 \times 10^{10} \) cm$^{-2}$, a gas density \( n_{\text{H}_2} \) from \( 1 \times 10^6 \) to \( 2 \times 10^8 \) cm$^{-3}$, and a temperature \( T \) from 80 to 200 K. We then simultaneously fitted the measured CH$_3$OH-A and CH$_3$OH-E line intensities via comparison with those simulated by the LVG model, leaving \( N_{\text{CH$_3$OH}}, n_{\text{H}_2}, T \), and the emitting size \( \theta \) as free parameters. Following the observations, we assumed the linewidths equal to 3.5 km s$^{-1}$ and 3.0 km s$^{-1}$ for 4A1 and 4A2, respectively, and we included the calibration uncertainty (15%) in the observed intensities.

The best fit is obtained for a total CH$_3$OH column density \( N_{\text{CH$_3$OH}} = 2.8 \times 10^{19} \) cm$^{-2}$ with reduced chi-square \( \chi^2 \) = 0.6 for 4A1 and \( N_{\text{CH$_3$OH}} = 1 \times 10^{19} \) cm$^{-2}$ with \( \chi^2 \rangle = 0.1 \) for 4A2. All the observed lines are predicted to be optically thick and emitted by a source of \( 0\arcsec22 \) for 4A1 and \( 0\arcsec24 \) for 4A2 (\( \sim 70 \) au) in diameter. Solutions with \( N_{\text{CH$_3$OH}} \geq 1 \times 10^{18} \) cm$^{-2}$ for 4A2 and \( \geq 1 \times 10^{19} \) cm$^{-2}$ for 4A1 are within 1\( \sigma \) of confidence level. Increasing the methanol column density, the \( \chi^2 \) decreases until it reaches a constant value for \( N_{\text{CH$_3$OH}} \geq 1 \times 10^{19} \) cm$^{-2}$ for 4A1 and \( N_{\text{CH$_3$OH}} \geq 3 \times 10^{19} \) cm$^{-2}$ for 4A2; this is because all the observed lines become optically thick (\( \tau \sim 10-30 \) for 4A1, \( \tau \sim 2-6 \) for 4A2) and, consequently, the emission becomes that of a blackbody. The results do not change assuming a linewidth \( \pm 0.5 \) km s$^{-1}$ with respect to the chosen one.

Figure 3 shows, for both sources, the density-temperature \( \chi^2 \) surface of the \( N_{\text{CH$_3$OH}} \) best fit. The gas temperature is (90–130) K for 4A1 and (120–190) K for 4A2, while for the gas density we obtained a lower limit of \( 2 \times 10^6 \) cm$^{-3}$ for 4A1 and \( 1.5 \times 10^7 \) cm$^{-3}$ for 4A2, which implies that the levels are LTE populated.

The fit results are reported in Table 2. The derived \( n_{\text{H}_2} \) and \( T \) are consistent with those computed with the model summarized in Su et al. \citeyear{Su2019} using our sizes.
from the frequency at 143 GHz toward 4A1 is 0.3, which leads to underestimate the predicted intensity with that observed by Taquet et al., assuming the best-achieving the predicted intensity with that observed by Taquet et al.

most stringent constraint to the dust optical depth, to compare the predicted intensity with that observed by Taquet et al. (2015). In the comparison, we took into account our LVG-derived source size and the angular resolution of the Taquet et al. (2015) observations. While for 4A2 we considered the line intensity quoted by Taquet et al. (2015), for 4A1, not having CH₃OH detection, we used the 3σ level of the Taquet et al. (2015) observations integrated over 3 km s⁻¹ (average linewidth toward 4A1; see Section 3).

The 4A1 and 4A2 CH₃OH predicted and observed values are reported in Table 2. While the two intensities are similar toward 4A2, they differ by about a factor five toward 4A1.

4.3. Dust Absorption Toward 4A1 and 4A2

Assuming that the difference between the predicted and observed intensities is due to the (foreground) dust absorption and using Equation (1), we derived the dust optical depth at 143 GHz (τ_{dust}^{143GHz}, Table 2). While τ_{dust}^{143GHz} toward 4A2 is small (~0.3), that toward 4A1 is large (>1.6) enough to attenuate the methanol line intensity by a factor >3. Therefore, the dust is affecting the millimeter line emission differently in the two sources.

5. Discussion

5.1. Is IRAS 4A1 a Hot Corino?

So far, only about a dozen hot corinos have been detected (Section 1) and the question arises whether this is because they are rare or because the searches have always been carried out at millimeter wavelengths, where dust could heavily absorb the line emission.

Our first result is that a source that was not supposed to be a hot corino based on millimeter observations, IRAS 4A1 (López-Sepulcre et al. 2017), indeed possesses a region with temperature ≥100 K (Section 3), namely the icy mantle sublimation one, and shows methanol emission (Section 3), the simplest of the iCOMs, when observed at centimeter wavelengths. According to its definition (Ceccarelli 2004), thus, IRAS 4A1 is a hot corino.

Although we cannot affirm that hot corinos are ubiquitous, it is clear that the searches at millimeter wavelengths may be heavily biased and that complementary centimeter observations are necessary to account for dust opacity and understand the occurrence of hot corinos.

5.2. 4A2 versus 4A1: Are They Chemically Different?

Unlike 4A2, no sign of iCOM millimeter emission was revealed toward 4A1 (Taquet et al. 2015; López-Sepulcre et al. 2017). Using ALMA observations at 250 GHz, López-Sepulcre et al. (2017) found that the iCOMs abundances toward 4A2 and 4A1 differ by more than a factor 17, with the largest values (~100) for HCOOCH₃ and CH₃CN.

The first question to answer is whether the chemical difference between the two coeval objects is real or due to a different absorption by the surrounding dust.

In Section 4.3, we found that τ_{dust} at 143 GHz toward 4A1 and 4A2 is ≳1.6 and 0.3, respectively (see Table 2). Using the dependence of τ_{dust} from the frequency (τ_{dust}/τ_{dust}^{100GHz} = (ν/ν_{dust})^{β}) and assuming β = 2 (ISM value), the optical depth scaled at 250 GHz (frequency at which López-Sepulcre et al. 2017, derived the above iCOMs abundance ratios) is ≳4.9 for 4A1 and 0.9 for 4A2. Therefore, the different dust absorption toward 4A1 and 4A2 provides us, as a lower limit, a factor of 55 difference in their line intensities (I₄/²/I₄/¹), comparable to the 4A2/4A1 iCOMs abundance ratios derived by López-Sepulcre et al. (2017). A large dust absorption was also suggested by the anomalous flattened continuum spectral index at 100–230 GHz (Li et al. 2017) and the 90° flipping of the linear polarization position angles observed at above and below 100 GHz frequencies (Ko et al. 2020).

Although we cannot exclude that a real chemical difference exists between 4A1 and 4A2, the observations so far available cannot support that hypothesis. Centimeter observations of other iCOMs than methanol are necessary to settle this issue. This conclusion may apply to other binary systems where an apparent chemical difference is observed using millimeter observations.

5.3. Are the iCOM Abundances in Hot Corinos Underestimated?

The dust absorption also affects the iCOM line intensities in 4A2. At 143 GHz τ_{dust} is 0.3, which leads to underestimate the iCOM abundances by about 30%. At higher frequencies, this factor becomes more important; e.g., at 250 GHz, where several hot corinos studies are carried out (see references in Section 1), the absorption factor would be 2.5, and at 350 GHz, frequency where the most sensitive iCOMs search has been carried out (e.g., Jørgensen et al. 2016), the absorption factor would be 6. This behavior also agrees with what has already been found in massive hot cores (e.g., Rivilla et al. 2017). Therefore, in order to derive reliable iCOM abundances, complementary centimeter observations are needed to estimate the dust absorption.
6. Conclusion

We carried out observations of methanol lines at centimeter wavelengths with the VLA interferometer toward the binary system IRAS 4A, where previous millimeter observations showed a possible chemical differentiation between the two objects. Specifically, while 4A2 showed iCOM line emission, 4A1 did not.

Our new observations detected 10 methanol lines in 4A1 and 4A2 with similar intensities. Using a non-LTE analysis and comparing with previous methanol millimeter observations, we showed that (1) 4A1 is a hot corino, (2) the lack of iCOMs detection toward 4A1 at millimeter wavelengths is caused by a large dust optical depth, and (3) the determination of the iCOMs abundances toward 4A2 via millimeter observations is slightly underestimated by the dust absorption.

Therefore, the difficulty in discovering new hot corinos could be because the searches have been carried out at (sub) millimeter wavelengths, where the dust absorption might not be negligible. The suspected different chemical nature of coeval objects of the same binary system needs also to be verified at centimeter wavelengths, as well as the iCOM abundances estimated from millimeter observations.

We conclude that centimeter observations of hot corinos are of paramount importance for their correct study. In the future, next generation instruments in the centimeter wavelength regime, such as ngVLA (McGuire et al. 2018) and SKA (Codella et al. 2015), could be the most efficient way to identify hot corinos and certainly the most appropriate facilities to study them.

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