Sub-mm gravity: confronting the modified dynamics with higher-dimensional theories

S.O. Mendes, R. Opher

Instituto Astronômico e Geofísico, Universidade de São Paulo,
Av. Miguel Stefano 4200, São Paulo, 04301-904 SP, Brazil

Abstract

We propose that future experiments aiming at the detection of deviations from the $1/r^2$ gravitational law on submillimetric scales can be used to test the modified Newtonian dynamics theory (MOND). Current experiments are able to test the gravitational field of masses $m \approx 1g$ at distances $r \approx 200 \mu m$, implying that they are probing accelerations well above the MOND limit ($a_0 \approx 1.2 \times 10^{-8} \text{cms}^{-2}$). We show that MONDian effects begin to be important at the submillimetric level for masses $m \leq 1 \text{mg}$. MOND makes predictions that are clearly distinguishable from those expected in a scenario with compact extra dimensions. This will enable direct confrontation between the two theories if future experiments can improve their mass scales to the milligram level.

Key words:
Gravitation: phenomenology, Gravitation: experimental tests
PACS: 04.80.Cc, 04.50.+h, 95.30.Sf

1 Introduction

There exists today an international effort aiming at the detection of deviations from conventional gravity, motivated by the possible existence of new spatial dimensions. The hypothesis of the existence of new spatial dimensions has been proposed as a solution to the hierarchy problem in particle physics. Compact dimensions would make the gravitational interaction appear to be weaker than other interactions, since part of its strength would “leak” into other dimensions [1]. Thus, the gravitational attraction should be stronger on small scales. Some

Email addresses: smendes@iagusp.usp.br (S.O. Mendes), opher@orion.iagusp.usp.br (R. Opher).
recent gravitational experiments have been searching for exactly this kind of deviation from Newtonian gravity [2]. We propose in this letter that the same experiments can be used to test for a similar prediction of deviations, suggested by the modified Newtonian dynamics (MOND), which is postulated to be valid for accelerations below a certain threshold, \( a_0 \approx 1.2 \times 10^{-8} \text{ cms}^{-2} \).

We briefly discuss the predicted accelerations for a given mass according to each theory, and then calculate the mass range which makes a direct comparison possible. We also discuss how MOND, in its original form, implicitly violates the strong equivalence principle and suggest that this hypothesis be relaxed. As we shall see, this is a necessary condition for testing MOND effects in the laboratory.

2 Theories with Extra Dimensions

The induced Yukawa-type gravitational potential in the context of extra dimensions can be written as [2,3]:

\[
V(r) = -\frac{GM}{r}(1 + \alpha e^{-r/\lambda}),
\]

where \( \alpha \) and \( \lambda \) are the intensity and range of the potential, respectively. According to recent experiments, the upper limit for \( \lambda \) is of the order of 1 mm. A review of theoretical and experimental constraints for both \( \lambda \) and \( \alpha \) can be found in [3]. Assuming spherical symmetry, the gravitational acceleration corresponding to the above potential is, thus,

\[
g_{\text{submm}} = \frac{GM}{r} \left[ \frac{1}{r} + \alpha e^{-r/\lambda} \left( \frac{1}{r} + \frac{1}{\lambda} \right) \right].
\]

3 Modified Newtonian Dynamics

The dark matter problem on galactic and extragalactic scales has led to the development of MOND [4]. This theory postulates that for accelerations below a certain threshold (determined empirically to be \( \approx 1.2 \times 10^8 \text{ cm s}^{-2} \)), Newtonian dynamics should be no longer valid. It was proposed, instead, that the correct form of the gravitational acceleration be given by

\[
g = \frac{g_N}{\mu(g/a_0)},
\]
where \( g_N \) is the usual Newtonian acceleration and \( \mu(x) \) is a function which obeys the relation

\[
\mu(x) = \begin{cases} 
1, & x \gg 1 \\
x, & x \ll 1 
\end{cases} 
\]  

(4)

A commonly used function having the required asymptotic behavior is [4]

\[
\mu(x) = \frac{x}{\sqrt{1 + x^2}}. 
\]  

(5)

Unfortunately, MOND cannot be considered to be a complete theory, since a relativistic theory whose weak field limit yields MOND does not exist. Some attempts to find a general gravitational theory have been made [5], but none of them were completely satisfactory. However, MOND fits rotation curve data from spiral galaxies very well with only one free parameter, while dark matter fits usually demand three free parameters (see [6] for a discussion). MOND also predicts the observed Tully-Fisher relation [7] for spiral galaxies (e.g. [6,8]), as well as the Faber-Jackson [9] relation for ellipticals, as long as M/L does not vary much with mass [10]. The theory has also been applied to other astrophysical phenomena, such as the stability and warp of disk galaxies [11], the internal structure of satellite galaxies [12], the fundamental plane of elliptical galaxies [13], and structure formation [14].

In addition to the absence of a relativistic theory which incorporates MOND, there is yet another difficulty to be dealt with. MOND does not appear to obey the strong equivalence principle (SEP). Milgrom [4] suggested, in a seminal paper, that the dynamics of a subsystem \( s \) of a system \( S \) should not be described by MOND when \( a_s \ll a_0 \) and \( a_S \gg a_0 \) (where \( a_s \) and \( a_S \) are the typical accelerations in \( s \) and \( S \), respectively). \( S \) and \( s \) could be, for example, a cluster of galaxies and a single galaxy belonging to this cluster, respectively. Thus, when \( a_s \ll a_0 \) and \( a_S \ll a_0 \), the accelerations are given by MOND, whereas when \( a_s \ll a_0 \) and \( a_S \gg a_0 \) they are not given by MOND, but rather by conventional Newtonian theory. This means that an observer inside an elevator in free fall, which is embedded in an external homogeneous gravitational field, would be able to detect this field. This obviously is a violation of the SEP and, if correct, would rule out any attempt of measuring MOND effects in the laboratory. Even if one could isolate a system where accelerations are well below \( a_0 \), the strong external fields from the Earth, Moon, and the Sun would erase any MOND signature, according to Milgrom [4].

The main motivation for introducing such a violation of the SEP apparently comes from the interpretation of data on open stellar clusters in the Solar Neighborhood. According to Jones [15], the dynamically deduced masses of
the clusters Pleiades and Praesepe are about 1.5 times as large as the mass that can be accounted for by the stars in these clusters. The fact that: (1) the internal accelerations within these clusters are a few times smaller than \(a_0\) (which would lead to a much larger mass discrepancy according to MOND); and (2) the accelerations of the clusters produced by the Galaxy is of the order of \(a_0\), led Milgrom to propose that MOND does not obey the SEP. A sample of only two open clusters seems to be too small to be used as a firm indication that MOND violates the SEP. A larger sample of both isolated and non-isolated star clusters could set the stage for a more complete study of the issue. More recent studies on wide binaries [16] and open clusters [17] also yield dynamical masses consistent with the stellar content. Newer and more refined data might help in this sense (e.g. [18]). In this work we neglect the problem of the violation of the SEP in MOND and examine the possibility of testing relations (3) and (5) in the laboratory.

### 4 MOND × Extra Dimensions

Test particles with masses small enough to make MONDian effects noticeable on submillimetric scales need to be studied. With the aid of eq. (5) and taking \(g_N\) as \(GM/r^2\), assuming spherical symmetry, we can solve eq. (3) for \(g\):

\[
g = \frac{1}{r^2} \left[ \frac{GM}{2} (GM + \sqrt{G^2M^2 + 4r^4a_0^2}) \right]^{1/2}.
\]

(6)

In order to estimate the limiting mass for which MOND predicts deviations from the \(1/r^2\) law, we must find the mass \(M\) which satisfies

\[
g = a_0.
\]

(7)

For example, \(M \approx 1\) miligram when \(r = 1\) mm. Current experiments on sub-mm gravity are able to test the gravitational field produced by test-particles with \(M \approx 1\) gram. Fig. (1) illustrates the type of deviations expected from both MOND and theories with extra dimensions. We used the \((\lambda - \alpha)\) diagram of [3] for eq. (2) and assumed that \(\lambda \approx 1\) mm and \(\alpha \approx 7\). If future experiments succeed in testing the required MOND mass scale at the sub-mm level, we will then be able to test for the presence of MONDian effects in strong external gravitational fields. MOND makes predictions that are clearly distinguishable from those of theories with new compact dimensions, which makes them competitive in interpreting future experimental results. The major problem is, of course, that very small masses are required for a direct comparison. We leave the question of whether such an achievement is plausible in the near future to the experimentalists. MOND could also be tested in other regions of the
mass-distance parameter-space: (60 g, 1 cm), (1 kg, 1 m), etc. In this case, one would test MOND alone, since this is well beyond the range expected for the strengthening of gravity due to extra dimensions.

Acknowledgements

The authors are grateful to C.D. Hoyle who kindly clarified some aspects of the experiment discussed in [2]. SOM would like to thank the Brazilian agency CAPES for financial support. RO would like to thank the Brazilian agencies FAPESP, CNPq and PRONEX/FINEP (no. 41.96.0908.00) for partial support. The authors also thank FAPESP for partial support through project no. 2000/06770-2.

References

[1] N. Arkani-Hamed, S. Dimopoulos and G. Dvali, Phys. Rev. D59, 086004 (1999).
[2] C. D. Hoyle, U. Schmidt, B. R. Heckel, E. G. Adelberger, J. H. Gundlach, D. J. Kapner and H. E. Swanson, Phys. Rev. Lett. 86, 1418 (2001).
[3] A. Kehagias and K. Sfetsos, Phys. Lett. B472, 39 (2000).
[4] M. Milgrom, ApJ 270, 365 (1983).
[5] J. D. Bekenstein and M. Milgrom, ApJ 286, 7 (1984); J. D. Bekenstein, in Second Canadian Conference on General Relativity and Relativistic Astrophysics, eds. A. Coley, C. Dyer and T. Tupper, World Scientific, Singapore, p.487 (1988); R. H. Sanders, ApJ 80, 492 (1997).
[6] R. H. Sanders and M. A. W. Verheijen, ApJ 503, 97 (1998).
[7] R. B. Tully and J. R. Fisher, A&A 54, 661 (1977).
[8] R. H. Sanders, ApJ 473, 117 (1996).
[9] S. M. Faber and R. E. Jackson, ApJ 204, 668 (1976).
[10] M. Milgrom, ApJ 287, 571 (1984).
[11] R. Brada and M. Milgrom, ApJ 519, 590 (1999); R. Brada and M. Milgrom, ApJ 531, L21 (2000a).
[12] D. Müller and R. Opher, ApJ 540, 57 (2000); R. Brada and M. Milgrom, ApJ 541, 556 (2000b).
[13] R. H. Sanders, MNRAS, 313, 767 (2000).
[14] R. H. Sanders, MNRAS, 296, 1009 (1998); R. H. Sanders, ApJ (accepted) astro-ph/0011439 (2001).

[15] B. F. Jones, AJ 75, 563 (1970); B. F. Jones, AJ 76, 470 (1971).

[16] L. M. Close, H. B. Richer and D. R. Crabtree, AJ 100, 1968 (1990).

[17] P. J. T. Leonard and D. Merritt, ApJ, 339, 195 (1989).

[18] ESA, The Hipparcos and Tycho Catalogues, ESA SP-1200 (1997)
Fig. 1. Gravitational acceleration as a function of distance produced by a test particle of mass 1 mg according to Newtonian theory (solid), MOND (dotted) and a higher dimensional theory (dashed).