Controlling Spin-Correlated Radical Pairs with Donor–Acceptor Dyads: A New Concept to Generate Reduced Metal Complexes for More Efficient Photocatalysis

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Abstract: One-electron reduced metal complexes derived from photoactive ruthenium or iridium complexes are important intermediates for substrate activation steps in photoredox catalysis and for the photocatalytic generation of solar fuels. However, owing to the heavy atom effect, direct photochemical pathways to these key intermediates suffer from intrinsic efficiency problems resulting from rapid geminate recombination of radical pairs within the so-called solvent cage. In this study, we prepared and investigated molecular dyads capable of producing reduced metal complexes via an indirect pathway relying on a sequence of energy and electron transfer processes between a Ru complex and a covalently connected anthracene moiety. Our test reaction to establish the proof-of-concept is the photochemical reduction of ruthenium(tris)bipyridine by the ascorbate dianion as sacrificial donor in aqueous solution. The photochemical key step in the Ru-anthracene dyads is the reduction of a purely organic (anthracene) triplet excited state by the ascorbate dianion, yielding a spin-correlated radical pair whose (unproductive) recombination is strongly spin-forbidden. By carrying out detailed laser flash photolysis investigations, we provide clear evidence for the indirect reduced metal complex generation mechanism and show that this pathway can outperform the conventional direct metal complex photoreduction. The further optimization of our approach involving relatively simple molecular dyads might result in novel photocatalysts that convert substrates with unprecedented quantum yields.

Introduction

Visible-light photoredox catalysis as a versatile tool in synthetic organic chemistry has attracted increasing attention in recent years,[1–8] and research on solar fuels continues to explore the possibility of harvesting sunlight to drive chemical reactions.[9–13] About half of the chemical syntheses in these fields are reductive transformations, with the most prominent key intermediates for substrate activation being one-electron-reduced Ru and Ir metal complexes. However, the direct generation of these intermediates via photoreduction of excited triplet states is usually inefficient, which is due to the heavy atom effect. Hence, a substantial part of the photons absorbed by the catalyst cannot be used productively, adversely affecting the sustainability aspect connected with photocatalysis. The reduction of any triplet state by a sacrificial electron donor D initially gives a so-called spin-correlated radical pair (upper right corners in Schemes 1a and b), which is, as a result of spin conversion, born in the multiplicity of its photoexcited precursor (i.e., in a triplet state).[14–16] Immediate recombination of this triplet radical pair in the solvent cage is spin-forbidden. In direct consequence, cage escape on a subnanosecond timescale producing non-correlated “free” radicals with typical lifetimes in the microsecond to millisecond range can proceed with unit inherent efficiency (η). This case of highly efficient formation of free radicals (Scheme 1a) that can react productively with substrates has been observed for many organic molecules lacking heavy atoms.[17–19] The second case, in which a triplet of a heavy-atom-containing molecule (e.g., a photoactive Ru complex M) is reduced, is more complicated. As a result of spin-orbit coupling, heavy atoms drastically accelerate radical pair intersystem crossing (ISC), such that triplet ([M•–D•−]) and singlet radical pairs ([M•+D•−]) interconvert right after electron-transfer quenching.[14–16] Only the latter radical pair can undergo ultrafast back electron transfer (bet) also known as in-cage or geminate recombination.[14] This unproductive pathway (highlighted with red labels in Scheme 1b) rapidly converts the absorbed photons into heat rather than generating reactive open-shell species.

Most photoredox studies merely analyze Stern–Volmer emission quenching to provide evidence for photoinduced electron
transfer processes. However, in photochemistry one excited state quenched via photoinduced electron transfer (PET) does not necessarily mean that one reduced or oxidized species is formed. Despite practically quantitative excited-state quenching, the actual yield of free radicals or radical ions might be close to zero, but the determination of cage-escape efficiencies for a given PET event requires more sophisticated experimental techniques such as quantitative transient absorption spectroscopy. Owing to the lack of systematic quantitative studies in photochemistry, the exact factors that govern the overall efficiencies of PET processes are still poorly understood and a reliable prediction of the values cannot be made in advance. For instance, the inherent photoreduction efficiencies for an anionic ruthenium(II) complex by a series of dianionic electron donors differ by as much as a factor of 30. Most values for PET reactions with triplet-excited Ru complexes are in the range from 0.05 (5%) to 0.6 (60%), clearly indicating that unproductive in-cage recombination is a general energy-wasting problem.

In this work, we have prepared novel molecular dyads consisting of a ruthenium-(tris)bipyridine unit (M) covalently attached to an anthracene (A) moiety, and investigated their photochemical properties. Using quantitative laser flash photolysis (LFP) with optical detection of the intermediates, we will present an alternative mechanism for the generation of one-electron reduced metal complexes \( M^{2+} \) that is able to outperform the conventional direct photoreduction of metal complexes. The underlying key principle is that the triplet of the heavy-atom free anthracene chromophore is photoreduced and the desired species \( M^{2+} \) is subsequently generated by a thermal (intramolecular) electron transfer, which does not suffer from spin-dependent loss channels.

**Results and Discussion**

Quantitative studies on cage escape yields or inherent efficiencies for PET reactions require the determination of molar absorption coefficients of the quenching products. Reductive quenching of triplet-excited ruthenium(tris)bipyridine \( ^3\text{Ru}(\text{bpy})_3^{2+} \leftarrow \left( ^1\text{M} \right) \) in water to yield \( M^{2+} \) is probably the best-investigated photoreduction of a heavy-atom containing chromophore. Among the readily available sacrificial donors, the ascorbate dianion \( \text{Asc}^{2-} \) is the most promising candidate for the alternative photoreduction mechanism presented herein, because of the very low potential for its oxidation. The direct photoreduction of \( ^3\text{Ru}(\text{bpy})_3^{2+} \) by \( \text{Asc}^{2-} \) (Scheme 2a) proceeds with an \( \eta \) of about 0.4, this means that some 60% of all quenching events do not produce the desired reduced complex \( \text{Ru}(\text{bpy})_3^{2+}+\left( ^1\text{M} \right)^{2+} \), which is a versatile intermediate for the initiation of numerous photoredox reactions. An \( \eta \) increase is highly desirable since it should not only result in a significantly shorter irradiation time for a given photoreaction, but it might also increase the catalyst stability by reducing the number of unproductive photocatalyst excitation events that lead to photodecomposition.

Guided by numerous studies on metal complex-aromatic hydrocarbon dyads as well as the photoreduction investigations of naphthalene and pyrene triplets, the following strategy (Scheme 2b) provides an indirect access to \( M^{2+} \) and aims to increase \( \eta \) by avoiding the direct photoreduction of \( ^3\text{Ru}(\text{bpy})_3^{2+} \): (1) Selective excitation of M to yield \( ^1\text{M} \) after ultrafast ISC, (2) intramolecular triplet-triplet energy transfer (TTET) producing \( ^1\text{A} \), (3) ary radical anion generation via photoreduction of \( ^1\text{A} \) by \( \text{Asc}^{2-} \) (ET), (4) intramolecular electron transfer (iET) from \( ^1\text{A} \) to M yielding \( M^{2+} \). In principle a strongly related reaction sequence could occur with both unconnected chromophores, but the linkage of the metal complex and the aromatic hydrocarbon unit in a dyad greatly enhances the overall efficiency as slow diffusion between desired reaction

![Diagram](image.png)
partners, which would compete with photophysical deactivation or radical recombination, is superfluous in our molecular dyads. All key steps are thermodynamically feasible (Scheme 2c) and the specific design of our molecular dyads is discussed in the next section. Interestingly, our approach is based on a “ping-pong” effect of energy and electron transfer in a molecular donor–acceptor system, which might be important for other applications beyond those related to photocatalysis: The donor of energy transfer is the acceptor of electron transfer.

Molecular design and ground state properties of the dyads

Aiming to produce a reduced ruthenium complex in a dyad, the first component of the dyad is predefined. We selected anthracene as second component serving as energy acceptor and redox relay on the following grounds: (i) Previous studies on molecular dyads reported very efficient anthracene triplet formation via TTET upon visible-light excitation of ruthenium-(tris)diimine complexes. \([47–51]\) (ii) The lowest triplet in these dyads is purely anthracene-localized and lives for at least several tens of microseconds at room temperature, \([47, 49, 50]\) ensuring sufficient time for reductive quenching (reaction step 3 in Scheme 2b). (iii) Inherent PET efficiencies \((\eta)\) close to unity have already been observed for one-electron oxidations of anthracene-localized triplets. \([46, 49, 52, 53]\) Consequently, steps 1–3 in Scheme 2b/c are expected to proceed with very high efficiencies.

Most prior investigations on Ru(bpy)\(_{2}\)\(^{2+}\)-anthracene dyads employed flexible spacers between the two chromophores. \([48, 50, 51, 53]\) However, dyads with flexible bridging units may exist in several different conformations including conformers with very short heavy atom (Ru)-anthracene distances. To avoid this potentially detrimental situation, which might promote rapid radical pair ISC, \([49]\) we employed \(p\)-xylene spacers (see Scheme 2b) to ensure rigid structures with well-defined interchromophore distances. According to our DFT calculations (see Section S3 of the Supporting Information) the Ru-anthracene distance is 12.6 Å in our dyad with one bridging unit (Ru(bpy))\(_{2}\)^{2+}-\(xy\)-Ant, whereas that distance amounts to 16.9 Å in the dyad with two \(p\)-xylene spacers (Ru(bpy))\(_{2}\)^{2+}-\(xy\)-Ant.

The dyads have been converted into their readily water-soluble nitrates by adding aq. KNO\(_3\) to the eluent mixture during the chromatographic purification of the final complexes. Synthetic procedures and characterization data can be found in Section S2 of the Supporting Information.

The UV/Vis absorption spectra of both dyads in aqueous solution are displayed in the upper panel of Figure 1. A comparison of these spectra with those of the reference compounds—Ru(bpy)\(_{3}\)^{3+} (blue line in Figure 1) and anthracene-9-propionate (purple line in Figure 1), which is a water-soluble anthracene derivative—clearly indicates the presence of both chromophores in the two dyads. The dyad spectra are basically a superposition of the individual Ru(bpy)\(_{3}\)^{3+} and anthracene spectra. This is due to the nature of the \(p\)-xylene spacers, as they do not significantly affect the electronic properties of covalently attached chromophores. \([54–57]\) In line with the expected poor orbital overlap, our DFT calculations (Supporting Information, Section S3) predict average equilibrium torsion angles between the \(p\)-xylene planes and the adjacent \(\pi\)-systems of the two chromophores on the order of 70°. Moreover, the computed
frontier orbitals (Supporting Information, Figure S2) are either purely Ru(bpy)$_3^{2+}$ (LUMO) or anthracene (HOMO) localized, without any noticeable orbital coefficients on the p-xylene bridging units. A closer look at the MLCT absorption bands (inset of Figure 1) reveals a minor red-shift (about 5 nm) for the dyads compared to conventional Ru(bpy)$_3^{2+}$ as well as very similar molar absorption coefficients. We, therefore, assume that the excited-state energies and redox potentials of the individual chromophores do not significantly alter in our dyads, and we took the pertinent literature values for Ru(bpy)$_3^{2+}$ and anthracene (see energy diagram in Scheme 2c). All steady-state (Figure 1) and time-resolved (Figure 2 and Figure 3) absorption investigations of this study were carried out in alkaline MilliQ water (50 mM NaOH, pH ~ 12.7). Working under these conditions ensures the presence of our electron donor ascorbate in its highly reactive dianionic form Asc$_2^@$ ($pK_a$ = 11.4,$^{46}$) more than 95% of Asc$_2^@$ in the acid-base equilibrium at pH 12.7), and does not negatively influence our photoredox systems shown in Scheme 2 in any way.

### Intramolecular triplet energy transfer

Selective excitation of the red edge of the MLCT absorption band$^{56}$ of both dyads with green (532 nm) laser pulses is feasible, as the UV/Vis spectra in Figure 1 clearly show. We expect the $^1$MLCT-excited Ru complex in the dyads to undergo ultra-fast$^{59}$ and quantitative$^{60}$ ISC to give a $^3$MLCT state, similar to what has been observed for free Ru(bpy)$_3^{2+}$, because compet-
ing energy (FRET) or electron transfer mechanisms with the co-
valently connected anthracene moiety can be ruled out for
thermodynamic reasons. Indeed, very weak \( ^1\text{MLCT} \) emission was observed during steady-state luminescence spectroscopy with deoxygenated \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) solutions, but a com-
parison with the parent compound \( \text{Ru(bpy)}_3^{2+} \) under identical excitation and detection conditions demonstrates that the dyad is less emissive by about two orders of magnitude. Initial time-resolved emission experiments on \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) upon excitation with 532 nm laser pulses of \( \sim 10 \text{ ns} \) duration showed an instrumentally limited decay of the \( ^1\text{MLCT} \) emission (detection wavelength, 600 nm). These results provide unam-
biguously fast and efficient \( ^1\text{MLCT} \) quenching in our dyad, given that the unquenched \( ^1\text{MLCT} \) lifetime of free \( \text{Ru(bpy)}_3^{2+} \) is about 600 ns under these experimental condi-
tions.\(^{37,61}\) To characterize the quenching process and the re-
sulting product, we carried out transient absorption investiga-
tions on short (Figure 2a) and long (Figure 2b) timescales. The experimental section with pertinent descriptions of the em-
ployed instruments as well as the underlying methods is given in the Supporting Information (Section S1.2).

Figure 2a displays picosecond transient absorption studies of deoxygcnated \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) and \( \text{Ru(bpy)}_3^{2+} \) \( ^1\text{MLCT} \) solutions at room temperature, which were excited with \( \sim 30 \text{ ps} \) pulses at 455 nm. We first investigated \( \text{Ru(bpy)}_3^{2+} \) \( ^1\text{MLCT} \) and reproduced the well-known transient absorption spectrum of its \( ^1\text{MLCT} \) state \( (\text{Ru(bpy)}_3^{2+}) \) with a pronounced absorption band at \( \sim 370 \text{ nm} \) and the characteristic ground state bleach in the blue region of the visible spectrum (blue spectrum in Fig-
ure 2a); both absorption features are produced instantaneously with our time resolution and remain persistent in the em-
ployed 5 ns time window.\(^{62}\) We speculated that the isosbestic point of \( \text{Ru(bpy)}_3^{2+} \) and \( \text{Ru(bpy)}_3^{2+} \) at 398 nm (compare, blue spectrum and kinetic trace in Figure 2a), which should also exist in the dyad, may be useful for studying the intramolecu-
lar quenching reaction (TTET from \( \text{Ru(bpy)}_3^{2+} \) to anthracene, step 2 in Scheme 2b, c) in isolation. Kinetic absorption measures-
ments with the dyad \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) show a clear first-
order signal rise at this main detection wavelength (detection window, 390–405 nm) with a time constant of about 130 ps (inset of Figure 2a). Practically identical time constants were observed at the \( ^1\text{MLCT} \) absorption maximum (360–375 nm) and in the spectral range of the \( \text{Ru(bpy)}_3^{2+} \) ground state bleach (450–475 nm).

The transient absorption spectrum of \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) obtained with a time delay ensuring quantitative \( \text{Ru(bpy)}_3^{2+} \) \( ^1\text{MLCT} \) quenching shows a clear (fine-structured) bleach in the near UV along with an intense absorption band with a maximum at 427 nm and a shoulder at 406 nm (Figure 2a, red spectrum), resembling the reported triplet-triplet absorption spectra of several anthracene derivatives.\(^{41,63-64}\) These findings are in per-
fact agreement with a purely anthracene-localized triplet state as TTET quenching product. The reverse TTET from the anthra-
cene triplet (triplet energy, \( \sim 1.84 \text{ eV} \)) to the ruthenium com-
plex (triplet energy, \( \sim 2.12 \text{ eV} \)) is endergonic by almost 0.3 eV, hence this reaction does not play any noticeable role in our dyad. Sub-100 ps (forward) energy transfer kinetics were reported for related ruthenium complex-anthracene dyads.\(^{51,67}\) In these dyads, the two chromophores were either directly con-
ected\(^{57}\) or a short flexible spacer was used. A longer inter-
chromophore distance and less pronounced orbital overlap can be expected for \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \), which might explain the slightly lower energy transfer rate that we observed in Figure 2a (127 ps). However, TTET in \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) is faster by about two orders of magnitude compared to rhenium(l) complex-anthracene dyads, probably due to inverted driving force effects for the latter.\(^{66}\)

After switching to our nanosecond laser setup (\( \sim 10 \text{ ns} \) laser pulse duration), we observed the very same transient absorption features of the sensitized anthracene triplet (Figure 2b). The formation of this species cannot be resolved with that instru-
ment, but it allows us to monitor the decay of the anthra-
cene-localized triplet. All absorption bands decay with very similar kinetics and return completely to the baseline on a 1 ms timescale, as is evident from the main plot of Figure 2b. Kinetic absorption traces (inset of Figure 2b) do not obey clear first-order kinetics. There is a slight admixture of second-order decay kinetics and we attribute that effect to the well-known annihilation of long-lived anthracene triplets. Time-gated emis-
sion spectra did not show delayed emission signals resulting from the annihilation product, i.e., singlet-excited anthracene, which is most likely due to intramolecular singlet energy trans-
fer in the dyad (FRET from singlet-excited anthracene to \( \text{Ru(bpy)}_3^{2+} \)).\(^{59}\) The natural (unquenched) triplet anthracene lifetime is on the order of 150 \( \mu \text{s} \) leaving ample time for bimo-
lecular reactions with suitable electron donors.

The \( ^1\text{MLCT} \) lifetime decrease from \( \sim 600 \text{ ns} \ (\text{Ru(bpy)}_3^{2+}) \) to 
\( \sim 130 \text{ ps} \) (dyad) indicates quantitative \( ^1\text{MLCT} \) quenching (99.98\%) in our dyad and based on our results, intramolecular TTET is the only quenching pathway. Therefore, it seems natu-
ral to assume a triplet anthracene formation quantum yield of 100\% upon visible excitation of \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \), and we per-
formed relative actinometry with \( \text{Ru(bpy)}_3^{2+} \) in water as widely used reference system,\(^{29,71,72}\) following the methodolo-
ger that we recently explained in detail.\(^{57}\) Using 532 nm laser pulses for excitation, a double determination of the molar ab-
sorption coefficient of \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) gave a value of \( 23000(\pm 500) \text{ M}^{-1} \text{ cm}^{-1} \) at the triplet anthracene absorption maximum, 427 nm. This molar absorption coefficient is compa-
rable to those determined for other triplets of substituted an-
thracene derivatives\(^{45,73}\) which substantiates our interpreta-
tion of a quantitative TTET without additional loss channels.

Similar nanosecond LFP investigations with the longer dyad \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) provide clear evidence for the fast sensi-
tized triplet anthracene formation (TTET in Scheme 2b), but they gave unexpectedly weak triplet anthracene signals for this dyad. Under identical excitation conditions (dyad absorp-
tion at the excitation wavelength and laser pulse energy, see Supporting Information, Figure S4 for details), the triplet an-
thracene signals for \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) are less intense by as much as a factor of three, compared to those of the \( \text{Ru(bpy)}_3^{2+} \cdot \text{xy} \cdot \text{Ant} \) dyad (Figure 2). We regard it as very unlikely that the molar absorption coefficients of both anthracene triplets differ significantly, in particular as their spectral shapes.
are virtually identical (Figure S4). Aggregation induced quenching might provide an explanation. The surfactant-like structure of Ru(bpy)$_3^{2+}$-xy$_2$-Ant with the hydrophilic Ru(bpy)$_3^{2+}$ chromophore as potential head group and the long hydrophobic π-system could promote the formation of aggregates or micellar structures. Ru(bpy)$_3^{2+}$-xy$_2$-Ant is readily water-soluble at the typical concentrations of our experiments (∼30 μm), but foam formation was clearly observed during Ar purging of the dyad solutions. The latter observation is in line with detergent properties of the longer dyad and might explain the seemingly low triplet anthracene formation efficiency observed for that compound. Further results that substantiate aggregation phenomena with Ru(bpy)$_3^{2+}$-xy$_2$-Ant are discussed in the next section. The results of this section, however, demonstrate that (i) the visible-light driven sensitization of the anthracene triplet is feasible with the novel dyads and (ii) that the shorter dyad Ru(bpy)$_3^{2+}$-xy$_2$-Ant is a promising candidate for the mechanism of Scheme 2, since the visible-light driven formation of the long-lived anthracene-localized triplet is both fast and quantitative.

Reactive quenching and formation of reduced metal complexes

The ascorbate dianion Asc$^{2-}$ with its redox potential of $E_{1/2}$ (Asc$^-$/Asc$^{2-}$) = 0.015 V vs. NHE [74] is a very potent electron donor. This reactant is evidently capable of reducing MLCT-excited ruthenium complexes with essentially diffusion-controlled kinetics [29, 30]. Furthermore, Asc$^{2-}$ should be able to quench the anthracene triplet reductively (step 3 in Scheme 2b), since its triplet energy (1.84 eV) slightly surpasses the sum of the potentials for Asc$^{2-}$ oxidation (0.015 V) [74] and anthracene reduction (1.71 V vs. NHE). [45] We indeed observed efficient quenching of the anthracene-localized triplet in Ru(bpy)$_3^{2+}$-xy$_2$-Ant, as the blue trace in Figure 3a illustrates. The anthracene triplet lifetime is reduced from about 150 μs to less than 1 μs, indicating complete quenching (> 99%). A new species is formed with the same rate as the Ru(bpy)$_3^{2+}$-xy$_2$-Ant decay (green trace in Figure 3a). A Stern-Volmer analysis of the purely dynamic quenching yielded a rate constant of $3.6 \times 10^9$ M$^{-1}$ s$^{-1}$. The transient absorption spectrum right after the quenching process (inset of Figure 3a) does not show the absorption bands derived from the initial quenching product, i.e., the anthracene radical anion $A^-$ (maxima at ∼700 nm and ∼360 nm) [75, 76]. However, the spectroscopic signatures of the reduced metal complex $M^+$ (Ru(bpy)$_3^{2+}$) are clearly observable, with characteristic bipyridine radical anion bands at 505 nm and 360 nm (superimposed by the ascorbate radical anion) [77] as well as the bleach of the MLCT absorption band of the parent ruthenium(II) complex [30, 78]. These results can be rationalized by a fast intramolecular electron transfer from the anthracene radical anion to the metal complex unit of the dyad (step 4 in Scheme 2b). The aren radical anion is significantly higher in energy than $M^+$ (Scheme 2c), and we expect the intramolecular electron transfer between $A^-$ and $M$ to occur on a (sub)nanosecond timescale. [79] The direct observation of the anthracene radical anion would thus not even be possible with much higher concentrations of the reductive quencher Asc$^{2-}$, because the monomolecular $A^-$ decay is always faster than its bimolecular formation (i.e., detectable anthracene radical anion concentrations cannot be formed for kinetic reasons). Nevertheless, the transient absorption studies of Figure 3a provide clear evidence for reductive quenching of the anthracene triplet and the formation of the desired reduced ruthenium complex $M^-$ as ultimate photoproduc in the dyad.

Figure 3b compares the $M^-$ formation results obtained for Ru(bpy)$_3^{2+}$ (blue) and our shorter dyad Ru(bpy)$_3^{2+}$-xy$_2$-Ant employing Asc$^{2-}$ as reductive quencher. To obtain conditions allowing a quantitative interpretation of the results, we (i) have chosen complex/dyad concentrations such that both solutions have identical absorbances at our excitation wavelength (532 nm), (ii) used a constant laser pulse energy throughout this series of experiments and (iii) doubled the Asc$^{2-}$ concentration for the well-characterized Ru(bpy)$_3^{2+}$ reference system. Not only do these conditions ensure identical triplet excited state concentrations right after the green laser pulses, but they also permit quantitative triplet quenching by the reactant Asc$^{2-}$ in both cases. The transient absorption spectra of the main plot of Figure 3b were recorded after the completion of reductive triplet quenching (zones marked in gray in the inset). The spectral shapes of the absorption bands, their positions as well as their relative intensities are almost identical. Based on these observations, we make the reasonable assumption of essentially identical molar absorption coefficients for $M^-$ regardless of whether the reduced complex is derived from conventional Ru(bpy)$_3^{2+}$ or the dyad Ru(bpy)$_3^{2+}$-xy$_2$-Ant. Hence, $M^-$ formation is in fact more efficient when using the dyad (Figure 3b), and the comparison of the signal intensities directly yields the relative inherent photoreduction efficiencies $\eta$. Taking the relative transient absorption values at the minimum (∼430 nm) and at both maximum (∼505 and ∼355 nm), we find that $\eta$ for Ru(bpy)$_3^{2+}$-xy$_2$-Ant is higher than that for Ru(bpy)$_3^{2+}$ by 48(±14)%. With the averaged literature $\eta$ value for the Ru(bpy)$_3^{2+}$ photoreduction by the ascorbate dianion ($\eta$ = 0.43) [29, 30] as reference, the photoreduction efficiency of the novel system, in which an anthracene triplet is reduced, amounts to about 0.64.

In a second comparison experiment, we employed the calibrated triplet anthracene $A^-$ absorption band of the dyad as internal reference signal (23 000 M$^{-1}$ cm$^{-1}$ at 427 nm) and compared its intensity to that of the absorption band of the photoproduct $M^-$ in the green spectral region. For that, we took the difference molar absorption coefficient of Ru(bpy)$_3^{2+}$ ($M^-$) at 505 nm ($\Delta \varepsilon = 14 200$ M$^{-1}$ cm$^{-1}$) [29, 30, 40] and calculated the relative concentrations of the anthracene triplet (from the amplitude of the blue trace in Figure 3a) ($\eta = 0$ value) and the photoproduct $M^-$ in the dyad (from the constant $\Delta \varepsilon$ value after quenching in the green trace in Figure 3a). This procedure gave an $\eta$ of 0.52 from the ratio between the concentrations of $M^-$ and $A^-$. The averaged $\eta$ resulting from our two independent determination methods is thus 0.58 ± 0.08.

This finding of an increased $\eta$ in the dyad compared to free Ru(bpy)$_3^{2+}$ represents clear evidence that a sophisticated indirect reaction sequence with several steps can be more efficient
than a direct single-step process with a large driving force (compare, Scheme 2), reminiscent of the antenna system and the redox cofactor chain in bacterial photosynthesis. In principle, our indirect visible light driven mechanism for the formation of M⁺ should be able to occur with up to unit inherent efficiency. A special geometrical arrangement between the ascorbyl radical and the reduced dyad might provide an explanation for the observable η of 0.58 in our case (Figure 4). Directly after the triplet anthracene reduction by Asc²⁻, the corresponding triplet radical pair is formed, and as indicated above, an ultrafast intramolecular electron transfer generates the reduced metal complex in that spin-correlated radical pair. Owing to Coulombic interactions, both between Ru(bpy)₃²⁺/Asc²⁻ before the quenching event and between Ru(bpy)₃²⁻/Asc⁻ after quenching, a rather close proximity between the two unpaired electrons and the heavy atom is assumed (see Figure 4). Such an arrangement could accelerate radical pair ISC and subsequent unproductive rearrangement processes, which compete with the desired cage escape producing free radical species.

In order to test this hypothesis and to find a system with η values as high as for purely heavy-atom free systems, we carried out similar laser experiments as in Figure 3 with the longer dyad Ru(bpy)₃²⁻-xy-Ant (see Supporting Information, Figure S5). As a result of the longer bridging unit, the ascorbyl radical produced via triplet anthracene quenching is expected to be more separated from the cationic ruthenium complex moiety thereby facilitating cage escape. In addition to the already mentioned unexpected findings that point to aggregation phenomena between Ru(bpy)₃²⁻-xy-Ant dyad molecules, we observed static quenching of the dyad triplet by Asc²⁻ and a slight change in the UV/Vis spectrum of Ru(bpy)₃²⁻-xy-Ant upon Asc²⁻ addition (see Supporting Information, Figure S6); both observations are in line with dyad-quencher preaggregation. The formation of M⁻ in the longer dyad is still feasible with the ascorbate dianion as quencher, but the estimated η for this system is as low as ~0.25 (Supporting Information, Section S4). We conclude that the aggregation and static quenching phenomena observed for this dyad make this compound unsuitable for exploring our concept of indirect M⁺ formation further, and thus the insight gained from Ru(bpy)₃²⁻-xy-Ant is limited. However, comparative DFT calculations on the one-electron reduced species derived from the parent ruthenium complex and both dyads were carried out. The spin density distributions for all energy-minimized open-shell structures are very similar (Supporting Information, Figure S3), and therefore these calculations support our experimental observations that M⁺ is the final quenching product in all cases.

The quest for other water-soluble reductive quenchers that can reduce both ¹MLCT and ³Ant was unsuccessful, in particular as the low excited state reduction potential of the anthracene triplet (about 0.1 V vs. NHE) militates for unusually strong reductants. Related spectroscopic studies were carried out with oxidative quenching of ruthenium complex-anthracene dyads some decades ago, when the power of photoredox catalysis was still hidden. Our proof-of-concept study for a more efficient indirect reductive quenching pathway avoiding the inefficient direct photoreduction of heavy-metal containing chromophores indicates that the molecular dyad approach could be developed into a general concept for improving the inherent efficiencies of PET reactions. This would have far-reaching implications for photocatalytic applications. The anthracene-containing dyads exhibit limited photostability, especially in the presence of dissolved oxygen (see Supporting Information for details). The exploitation of molecular dyads comprising a well-established 4d or 5d metal complex and a purely organic chromophore for more efficient photocatalytic transformation thus requires the design of more robust dyad-quencher systems, and our study paves the way for this promising idea.

Finally, we compare the quenching efficiencies of triplet-excitated Ru(bpy)₃²⁻-xy-Ant with that of conventional Ru(bpy)₃²⁻. Using the well-known photokinetic equations together with the unquenched triplet lifetimes and the bimolecular rate constants for reductive quenching with Asc²⁻, we simulated the quenching efficiencies as a function of the Asc²⁻ concentrations (Figure 5). The long natural lifetime of the anthracene-localized triplet in Ru(bpy)₃²⁻-xy-Ant (red lines) clearly overcompensates the slower Asc²⁻ quenching rate constant, as is evident from the much higher quenching efficiencies for the dyad compared to the conventional complex at given quencher concentrations (compare, blue and red thin lines in Figure 5). The half-quenching concentration of Asc²⁻ is as low as 18.5 μM for the dyad, whereas that quantity is 260 μM for Ru(bpy)₃²⁻. After multiplication of the quenching efficiency curves with the respective inherent PET efficiencies, η, the concentration-dependent cage escape yields result (thick lines in Figure 5). These curves reflect the yields of reduced metal complexes M⁻ capable of activating suitable substrates (see also panels a and b in Scheme 2) or co-catalysts. The analysis of the low Asc²⁻ concentration range in Figure 5 unambiguously establishes that M⁻ formation can indeed be more efficient by an order of magnitude when using our novel dyad Ru(bpy)₃²⁻-xy-Ant.
Figure 5. Comparison of quenching (thin lines) and cage escape efficiencies (thick lines) for the photoreductions of $\text{Ru(bpy)}_{3}^{2+}$ to $\text{xy}_1$-Ant and $\text{Ru(bpy)}_{3}^{3+}$ triplet by $\text{Asc}^{2-}$ as a function of the donor ($\text{Asc}^{2-}$) concentration. Parameters used for the kinetic simulations (values for the dyad shown in brackets): unquenched triplet lifetime, 0.6 μs (150 μs); quenching rate constant, 6.4 × 10^{10} \text{ M}^{-1} \text{ s}^{-1} (1.26 × 10^{10} \text{ M}^{-1} \text{ s}^{-1}); inherent cage escape yield, 0.43 (0.58). See text for further details.

Conclusions

In summary, we have developed a novel mechanism for the indirect photochemical generation of one-electron reduced metal complexes $\text{M}^{n+}$ (see Figure 6 for the key steps and pertinent kinetic constants), which are key intermediates in photoredox catalysis but whose direct photochemical generation is inherently inefficient. Most studies on visible-light driven photoredox catalysis completely ignore the efficiency aspect (i.e., the overall quantum yields) and merely focus on the feasibility to use visible light as energy input. In addition to photoredox catalysis, the $\text{M}^{n+}$ generation is important for many photochemical carbon dioxide reduction\[29, 30\] and hydrogen production\[31-34\] mechanisms, demonstrating that the new mechanism introduced in this manuscript has several possible application areas. Our study on the heavily underexplored inherent efficiency of photoinduced electron transfer events might trigger further quantitative investigations that could contribute to a better understanding of how to use photons more efficiently,\[35\] which could ultimately result in more sustainability in photochemistry.\[36\] Elucidating the interplay of spin states, heavy atom effects and inherent reaction (rather than emission quenching) efficiencies should therefore be very important for the photochemistry of the future.\[37-39\] Novel molecular dyads with a visible-light-harvesting metal complex and a redox-active heavy-atom free chromophore promise a promising alternative to conventional metal complexes, as has emerged from this work.

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Conflict of interest

The authors declare no conflict of interest.

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