Landau parameters of nuclear matter in the spin and spin-isospin channels

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The equation of state (EOS) of nuclear matter is a major issue of the quantum-mechanical many-body problem. The reliability of any theoretical prediction has been measured on its capability of reproducing the empirical saturation energy and density, once the convergence of the theory has been firmly established. This is the case of the Brueckner-Bethe-Goldstone theory: the hole-line expansion has been proved in fact to be rapidly converging \cite{3}, and the Brueckner-Hartree-Fock (BHF) approximation implemented by three-body forces (3BF) can account for the empirical saturation point \cite{2}. Besides saturation density and energy additional constraints have been the compressibility extracted from the monopole energy and the symmetry energy from the binding energy of \( N \neq Z \) nuclei. In the spin-isospin channel a further constraint has been provided by the Gamow-Teller (GT) giant resonances (see \cite{5} for a review). Recently the Landau parameter \( G_0' \) has been extracted from the experimental data with a very small uncertainty \cite{6}. So it represents a very robust constraint for the EOS of nuclear matter. The theoretical prediction of \( G_0' \) demands for extending the calculation of the EOS of nuclear matter to spin-polarized neutrons and protons. Such calculations have been stimulated by the search of a spontaneous transition to a ferromagnetic state which can have important implication in the physics of atomic nuclei. We extended the calculation of the EOS to polarized nuclear matter in the framework of the BHF approximation with the continuous choice for the auxiliary potential. As 2BF we took the Argonne \( AV_{18} \) \cite{5} and as 3BF the one from Ref.\cite{6}. Starting from unpolarized nuclear matter at a given density \( \rho \) (we only consider symmetric nuclear matter), we ran different polarization states of neutrons and protons

\[
\delta_n = \frac{N_\uparrow - N_\downarrow}{N}, \quad \delta_p = \frac{Z_\uparrow - Z_\downarrow}{Z}, \quad \rho = \frac{N + Z}{V}. \tag{1}
\]

The results are reported in Fig. 1 for two typical densities of nuclear matter with 2BF (left side) and 2BF plus 3BF (right side). Since the EOS of isospin-polarized nuclear matter fulfills a quadratic law as a function of isospin-symmetry parameter \( \delta = \delta_n = \delta_p \), and the same is true for the EOS of spin-polarized nuclear matter versus the spin-symmetry parameter \( \delta_p \), we fit the EOS (data plotted as symbols in Fig.1) by the least square method in the mixed case according to a quadratic law

\[
E_A(\rho, \delta_n, \delta_p) - E_A(\rho, 0, 0) = \sum_{\tau \tau'} \Lambda_{\tau\tau'}(\rho) \delta_\tau \delta_{\tau'}, \tag{2}
\]

where \( \tau = n, p \) is the isospin quantum number. In symmetric nuclear matter (SNM) the coefficients \( \Lambda_{\tau\tau'} \) are related to the zero-order Landau parameters

\[
G_0 = G_{nn}^0 + G_{np}^0 = \frac{4N(0)}{\rho} (\Lambda_{nn} + \Lambda_{np}) - 1, \tag{3}
\]
\[
G_0' = G_{nn}^0 - G_{np}^0 = \frac{4N(0)}{\rho} (\Lambda_{nn} - \Lambda_{np}) - 1. \tag{4}
\]

The results are reported in Fig. 1 for two typical densities of nuclear matter with 2BF (left side) and 2BF plus 3BF (right side). Since the EOS of isospin-polarized nuclear

FIG. 1: EOS of spin-polarized nuclear matter. The symbols are from microscopic calculations, and the lines are drawn only to guide the eye.
where \( N(0) \) is the level density at the Fermi surface. In Fig. 2 the two Landau parameters are plotted as a function of the density based on Eqs. (3) and (4).

The Landau parameter \( G_0' \) is the strength of the spin-isospin component \( V_{\sigma T} = G_0' (\tau_1 \cdot \tau_2) (\sigma_1 \cdot \sigma_2) \) of the residual interaction \( I_2 \), which governs the Gamow-Teller (GT) giant resonance in nuclei. Its value at the saturation point has been determined with high precision from the experimental excitation energy of the GT resonance on \( ^{90}\text{Ni} \). The value reported is \( 1.182 < G_0' < 1.188 \) (we have multiply by a factor two according to the definition we adopted). More recent fit on \( ^{112}\text{Sn} \) and \( ^{208}\text{Pb} \) within a RPA calculation with Skyrme forces confirm such a prediction. Our nuclear-matter prediction of \( G_0' \) at the saturation density, which is about 1.22 including 3BF is in a pretty good agreement with the previous values. The value without 3BF of about 1.30 is in less agreement. Since the behaviour of \( G_0' \) around the saturation point is very flat, there is no room for large uncertainties in the comparison. Such an agreement provides a further support to the important role played by the microscopic 3BF as to reproducing all saturation properties of nuclear matter. Other predictions of \( G_0' \) including 3BF are from phenomenological Skyrme forces, which unfortunately are scattered in wide range of values lower than the experimental one.

So far experimental information on \( G_0 \) is not enough since spin resonances have only been observed with too small strength compared to other collective modes.

The prediction of \( G_0 \) and \( G_0' \) for densities other than the nuclear density, which is reported in Fig. 2, is of great interest in the study of neutron stars. In connection with the strong magnetic fields observed in neutron stars some authors \( 5, 6, 7, 14 \) studied the magnetic susceptibility \( \chi \) in neutron matter and found that \( G_0 \) reduces the \( \chi \) of degenerate neutron gas. This reduction is amplified at high density when including 3BF either in Brueckner calculations and in Montecarlo many-body simulations.

Spin and spin-isospin excitations of nuclear matter are coupled to the weak interaction governing the neutrino emission of URCA processes as well as the neutrino transport in neutron stars. The high-density increase of \( G_0 \) and \( G_0' \), driven by the 3BF, is expected to have important implications for the neutron star cooling for the sizeable enhancement induced by the nuclear correlations on the neutrino free path.

In conclusion, in this note we reported on a BHF calculation of the Landau parameters \( G_0 \) and \( G_0' \) as a function of baryonic density. The main scope was to point out the large effect of 3BF, especially at high densities. At the 2BF level, there is a wide disagreement with previous Brueckner calculations (see Ref. \( 10 \) and Refs. therein quoted) that has not yet clearly explained, since one can hardly control and compare the different approximations. On the other hand, our prediction for \( G_0' \) has been found to be in very good agreement with the experimental value extracted from GT resonance when 3BF is included. The relevance of the spin Landau parameter for neutron stars has been also discussed in connection with magnetic susceptibility and neutrino mean free path.

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[1] H.Q. Song, M. Baldo, G. Giansiracusa and U. Lombardo, Phys. Rev. Lett. 81, 1584 (1998).
[2] W. Zuo, A. Lejeune, U. Lombardo, J.-F. Mathiot, Nucl. Phys. A 706, 418 (2002).
[3] F. Osterfeld, Rev. Mod. Phys. 64, 491 (1992).
[4] T. Suzuki and H. Sakai, Phys. Lett. B 455, 25 (1999).
[5] I. Vidaña, A. Polls and A. Ramos, Phys. Rev. C 65, 035804 (2002).
[6] W. Zuo, U. Lombardo and C.W. Shen, in Quark-Gluon Plasma and Heavy Ion Collisions, edited by W.M. Alberico et al. (World Scientific, Singapore, 2002), p.192.
[7] I. Vidaña and I. Bombaci, Phys. Rev. C 66, 045801 (2002).
[8] R.B. Wiringa, V.G.J. Stocks, R. Schiavilla, Phys. Rev. C 51, 38 (1995).
[9] P. Grange, A. Lejeune, M. Martzolff, J.-F. Mathiot, Phys. Rev. C 40, 1040 (1989).
[10] W. Zuo, A. Lejeune, U. Lombardo and J.-F. Mathiot, Eur. Phys. Journ. A 14, 469 (2002).
[11] I. Bombaci and U. Lombardo, Phys. Rev. C 44, 1892.
(1991).

[12] A.B. Migdal, *Theory of finite Fermi systems and applications to atomic nuclei* (Interscience Publ., New York, 1967).

[13] M. Bender, J. Dobaczewski, J. Engel and W. Nazarewicz, Phys. Rev. C 65, 054322 (2002).

[14] S. Fantoni, A. Sarsa and K.E. Schmidt, Phys. Rev. Lett. 87, 181101 (2001).

[15] U. Lombardo, C.W. Shen, N. Van Giai and W. Zuo, in Proceedings of International Symposium on Physics of Unstable Nuclei (Halong Bay, Vietnam, Nov. 20-25, 2002), to appear in Nuclear Physics (May 2003).

[16] M. Baldo, L.S. Ferreira, Phys. Rev. C 50, 1887 (1994).