The Random Transiter - EPIC 249706694/HD 139139

Rappaport, S.; Vanderburg, A.; Kristiansen, M. H.; Omohundro, M. R.; Schwengeler, H. M.; Terentev, I. A.; Dai, F.; Masuda, K.; Jacobs, T. L.; LaCourse, D.

Published in:
Monthly Notices of the Royal Astronomical Society

Link to article, DOI:
10.1093/mnras/stz1772

Publication date:
2019

Document Version
Publisher's PDF, also known as Version of record

Link back to DTU Orbit

Citation (APA):
Rappaport, S., Vanderburg, A., Kristiansen, M. H., Omohundro, M. R., Schwengeler, H. M., Terentev, I. A., ... Mann, A. (2019). The Random Transiter - EPIC 249706694/HD 139139. Monthly Notices of the Royal Astronomical Society, 488(2), 2455-2465. https://doi.org/10.1093/mnras/stz1772

General rights
Copyright and moral rights for the publications made accessible in the public portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

- Users may download and print one copy of any publication from the public portal for the purpose of private study or research.
- You may not further distribute the material or use it for any profit-making activity or commercial gain
- You may freely distribute the URL identifying the publication in the public portal

If you believe that this document breaches copyright please contact us providing details, and we will remove access to the work immediately and investigate your claim.
The Random Transiter – EPIC 249706694/HD 139139

S. Rappaport,1 A. Vanderburg,2† M. H. Kristiansen,3,4 M. R. Omohundro,5 H. M. Schwengeler,5 I. A. Terentev,5 F. Dai,1,6 K. Masuda,6† T. L. Jacobs,7 D. LaCourse,8 D. W. Latham,9 A. Bieryla,9 C. L. Hedges,10 J. Dittmann,11 G. Barentsen,10 W. Cochran,2 M. Endl,2 J. M. Jenkins12 and A. Mann13

1 Department of Physics, and Kavli Institute for Astrophysics and Space Research, M.I.T., Cambridge, MA 02139, USA
2 Department of Astronomy, The University of Texas at Austin, 2515 Speedway, Stop C1400, Austin, TX 78712
3 DTU Space, National Space Institute, Technical University of Denmark, Elektrovej 327, DK-2800 Lyngby, Denmark
4 Brorfelde Observatory, Observator Gyldenkernes Vej 7, DK-4340 Tølløse, Denmark
5 Citizen Scientist, c/o Zooniverse, Department of Physics, University of Oxford, Denys Wilkinson Building, Keble Road, Oxford, OX1 3RH, UK
6 Department of Astrophysical Sciences, Princeton University, Princeton, NJ 08544, USA
7 Amateur Astronomer, 12812 SE 69th Place Bellevue, WA 98006, USA
8 Amateur Astronomer, 7507 52nd Place NE Marysville, WA 98270, USA
9 Center for Astrophysics | Harvard & Smithsonian 60 Garden Street, Cambridge, MA 02138, USA
10 Kepler/K2 Guest Observer Office, NASA Ames Research Center, CA 94035, USA
11 Department of Earth and Planetary Sciences and Kavli Institute for Astrophysics and Space Research, M.I.T., Cambridge, MA 02139, USA
12 NASA Ames Research Center, Moffett Field, CA 94035, USA
13 Department of Physics and Astronomy, University of North Carolina at Chapel Hill, Chapel Hill, NC 27599-3255, USA

Accepted 2019 June 24. Received 2019 June 22; in original form 2019 May 22

ABSTRACT

We have identified a star, EPIC 249706694 (HD 139139), that was observed during K2 Campaign 15 with the Kepler extended mission that appears to exhibit 28 transit-like events over the course of the 87-d observation. The unusual aspect of these dips, all but two of which have depths of 200 ± 80 ppm, is that they exhibit no periodicity, and their arrival times could just as well have been produced by a random number generator. We show that no more than four of the events can be part of a periodic sequence. We have done a number of data quality tests to ascertain that these dips are of astrophysical origin, and while we cannot be absolutely certain that this is so, they have all the hallmarks of astrophysical variability on one of two possible host stars (a likely bound pair) in the photometric aperture. We explore a number of ideas for the origin of these dips, including actual planet transits due to multiple or dust emitting planets, anomalously large TTVs, S- and P-type transits in binary systems, a collection of dust-emitting asteroids, ‘dipper-star’ activity, and short-lived starspots. All transit scenarios that we have been able to conjure up appear to fail, while the intrinsic stellar variability hypothesis would be novel and untested.

Key words: stars: activity – circumstellar matter – stars: general – planets and satellites: general.

1 INTRODUCTION

The K2 mission (Howell et al. 2014), spanning 4 yr, has yielded a large number of impressive results regarding what are now considered conventional areas of planetary and stellar astrophysics. Nearly a thousand new candidate and confirmed exoplanets have been reported (see e.g. Mayo et al. 2018; Dattilo et al. 2019). In addition, the K2 mission has continued the Kepler (Koch et al. 2010) discovery of dynamically interesting triple and quadruple star systems (e.g. Rappaport et al. 2017; Borkovits et al. 2019). Numerous asteroseismological studies have also been carried out with K2, including, e.g. Chaplin et al. (2015), Lund et al. (2016a, 2016b), and Stello et al. (2017).

In addition to these areas of investigation, K2 has taken some departures from conventional stellar and planetary astrophysics, and uncovered a number of unexpected treasures. These include periodic transits by dusty asteroids orbiting the white dwarf WD 1145 + 017 (Vanderburg et al. 2015); a disintegrating exoplanet in a 9-h orbit...
2.2 Search for unusual objects

In addition to using Box Least Squares (‘BLS’; Kovács, Zucker & Mazeh 2002), searches for periodic signals, most of us (MHK, MO, IT, HMS, TJ, DL) continued our visual survey of the K2 light curves searching for unusual objects. This has led to a number of discoveries not initially caught by BLS searches, including multiply eclipsing hierarchical stellar systems (e.g. EPIC 220209460, Rappaport et al. 2017; EPIC 219271635, Borkovits et al. 2018; EPIC 249432662, Borkovits et al. 2018); multiplanet systems (e.g. WASP 47, Becker et al. 2015; HIP 41378, Vanderburg et al. 2016; EPIC 248435473, Rodriguez et al. 2018); stellar eclipses by accretion discs (EPIC 225300403, Zhou et al. 2018); and a single very deep stellar occultation (80 per cent deep in EPIC 204376071, Rappaport et al. 2019).

Here, we present the discovery, via a visual survey, of more than two dozen shallow, transit-like events with durations of 0.75–7.4 h identified in the C15 light curve of EPIC 249706694 (aka HD 139139; GO-program PIs: Charbonneau, GO15009; Howard, GO15021; Buzasi, GO15034).3 The strange and intriguing thing about these transit-like events is that they exhibit no obvious periodicity. In fact, EPIC 249706694 was not flagged by any BLS searches of which we are aware, and only visually spotted while inspecting the light curves using LcTools4 (Kipping et al. 2015).

Though the transit-like events were first spotted in the Kepler-pipeline light curves (Jenkins et al. 2010), we utilized the data set of Vanderburg & Johnson (2014) to construct the light curve that is then used throughout the remainder of this study. The Vanderburg & Johnson (2014) pipeline generates light curves starting with pixel-level data. We then subjected this light curve to a high-pass filter to remove the starspots activity. The top panel of Fig. 1 shows the raw data set for 87 d of the K2 observation. The middle panel presents the data after high-pass filtering, and this more clearly reveals the 28 transit-like events that we have detected. The bottom panel zooms in on a 15-d segment of the data set around four of the transit-like events.

We tabulate all the depths and durations of the 28 dips in Table 1. The robustness of the light curve we used was checked by varying both the size and shape of the photometric aperture, and the self flat-fielding (‘SFF’) parameters, such as the spacing of knots in the spline that models low-frequency variations, the size of the windows on which the one-dimensional flat-fielding correction was performed, and the number of free parameters describing the flat field. The times, depths, and widths of the dips are quite invariant to these changes. We also downloaded a light curve of EPIC 249706694 generated with the EVEREST pipeline (Luger et al. 2016), and find excellent agreement with our data set.

2.3 Archival data on EPIC 249706694

In Table 2, we collect the various archival magnitudes as well as the Gaia kinematic data on EPIC 249706694 (HD 139139). The target star is quite bright at $K_p = 9.68$ and of solar effective temperature and radius. However, there is also a neighbouring star (hereafter ‘star B’) that is 3 mag fainter in $G$ somewhat cooler (4400 K), and separated from the target star by 3.3 arcsec. There is almost no additional Gaia information on this star, so we have no initially clear picture of whether star B is a bound companion to star A or not.

We have done our own astrometry on the elongated 2MASS image (from 1998.4, Skrutskie et al. 2006), the PanSTARRS image (from 2012.1 and saturated in places; Flewelling et al. 2016), and a current-epoch (2019.3) image we acquired at the Fred Lawrence Whipple Observatory (Mt. Hopkins) using the KeplerCam camera on the 48-inch telescope. Using these data we can typically find the difference in RA and Dec of stars A and B to accuracies of 0.4 arcsec, 0.1 arcsec, and 0.2 arcsec, for the three images, respectively. This is sufficiently good, over the ±21 yr baseline to conclude that the proper motion of star B in both directions tracks that of star A to within 20 mas yr$^{-1}$. Our astrometry is summarized in Table 3. Given that the proper motions in RA and Dec of star A

---

3https://keplerscience.arc.nasa.gov/k2-approved-programs.html#campaing
4https://sites.google.com/a/lctools.net/lctools/
Figure 1. K2 light curve for EPIC 249706694 (HD 139139) during Campaign 15. Top panel: The raw 87-d light curve. Middle panel: Light curve after filtering out the slow modulations due to star spots and trends that result from data processing. There are 28 transit-like events whose flux dips below the level of the red horizontal line, which we take to be significant (at 82 ppm below unity). Bottom panel: a shorter 15-d segment of the light curve containing four of the transit-like events.

Table 1. Transit-like events of the Random Transiter.

| Number | Time (BKJD<sup>a</sup>) | Duration<sup>a</sup> (h) | Depth<sup>a</sup> (ppm) |
|--------|-------------------------|-------------------------|-------------------------|
| 1      | 3161.6826(25)           | 1.14(17)                | 266(33)                 |
| 2<sup>c</sup> | 3163.3811(184)        | 4.41(1.25)              | 67(13)                  |
| 3      | 3164.7902(34)           | 2.72(23)                | 244(19)                 |
| 4      | 3165.6969(35)           | 2.33(25)                | 224(27)                 |
| 5      | 3166.0265(30)           | 1.08(17)                | 192(35)                 |
| 6      | 3166.3137(22)           | 1.10(11)                | 218(34)                 |
| 7      | 3174.3337(29)           | 2.04(15)                | 274(25)                 |
| 8      | 3175.3532(47)           | 1.75(26)                | 130(24)                 |
| 9      | 3178.0673(29)           | 2.55(19)                | 253(20)                 |
| 10     | 3181.4968(53)           | 4.22(36)                | 189(17)                 |
| 11     | 3186.5145(43)           | 2.09(25)                | 189(23)                 |
| 12     | 3190.5272(30)           | 2.20(19)                | 222(20)                 |
| 13     | 3195.8316(48)           | 1.33(28)                | 128(26)                 |
| 14     | 3196.8193(34)           | 2.73(23)                | 226(19)                 |
| 15     | 3202.4718(42)           | 4.79(28)                | 220(22)                 |
| 16     | 3202.9034(49)           | 1.67(36)                | 128(39)                 |
| 17     | 3205.6068(18)           | 2.77(12)                | 408(18)                 |
| 18     | 3209.7655(30)           | 1.14(19)                | 210(32)                 |
| 19     | 3216.4965(66)           | 8.19(45)                | 256(12)                 |
| 20     | 3221.1631(114)          | 0.95(63)                | 257(29)                 |
| 21     | 3223.6633(47)           | 4.21(33)                | 222(15)                 |
| 22     | 3225.7971(67)           | 7.40(47)                | 212(13)                 |
| 23     | 3226.5548(77)           | 5.92(53)                | 176(13)                 |
| 24     | 3231.0569(49)           | 5.41(34)                | 247(12)                 |
| 25     | 3238.5381(55)           | 1.17(26)                | 196(32)                 |
| 26     | 3240.6550(50)           | 6.42(34)                | 237(12)                 |
| 27     | 3240.8371(29)           | 1.07(23)                | 228(29)                 |
| 28     | 3243.5688(28)           | 0.74(14)                | 200(29)                 |

Notes. <sup>a</sup>The depths and durations were determined in two independent ways: (i) fitting to planet transit models and (ii) fitting modified hyperbolic secants (i.e. with a squared argument), and these yielded consistent results. <sup>b</sup>We define ‘BKJD’ as BJD-2454833. <sup>c</sup>This dip feature may be suspect because it is the most shallow and a K2 thruster firing occurs within it.

ANALYSIS OF THE DIPS

We have searched extensively for periodicities among the 28 transit-like events. In the top panel of Fig. 2, we show the box least-squares spectrum. There are no significant peaks. The middle panel of Fig. 2 is the result of a Lomb–Scargle transform, and, as can be expected if the BLS shows no signal for transit-like events, neither will an LS transform.

In the bottom panel of Fig. 2, we show the results of an ‘Interval Match Transform’ (‘IMT’; see section 7 of Gary et al. 2017) which consists of the brute force testing of 10<sup>5</sup> trial periods against all...
unique (28 × 27/2 = 378) dip arrival-time differences. For each trial period, we checked whether the time interval between each distinct pair of dips was equal to an integer number of the trial period. We allowed for a plus or minus 3 per cent leeway in terms of matching an integral number of cycles. (This tolerance for matching an integer can be set arbitrarily to allow for small or even large TTVs.) The IMT in this case should produce a 6 per cent accidental matches. Therefore, if even 1/4 of the 28 events are truly periodic, then there should be an additional 21 (i.e. 7 × 6/2) consequential matches at the true periodicity. And, we see no such indication.

From these tests, we know that not more than about 7 of the 28 dips in flux can be periodic. But, perhaps there are small-number sets of periodic dips. In order to test just how sensitive the BLS algorithm is to a limited set of periodic dips in the presence of numerous non-periodic dips, we did the following exercise. First, we injected a set of 10 periodic transits of depth 200 ppm and duration 3 h into the actual data set containing the 28 apparently random transits. This choice of depth and duration seems reasonably typical of the other dips. Those 10 periodic transits were easily detected at high signal to noise (S/N) with the BLS search. We then repeated the exercise, but with 9, 8, 7, ... etc. added periodic dips. They were readily detectable down to the addition of only four periodic dips. Three dips (and obviously any fewer) could not be convincingly detected. In Table 4, we list the number of added periodic transits to the data set, and the S/N with which they were detected. From this, we conclude that any periodicities among the 28 dips in flux can occur in sets no larger than four events each.

![Figure 2. Three transforms used to search for periodicities in the 28 transit-like events. Top panel: Box least squares transform reveals no significant periods in the range of 0.1–50 d. Middle panel: Lomb–Scargle transform also shows no significant periods over the same range as the BLS. Bottom panel: An 'Interval Match Transform' (see Gary et al. 2017), which allows for very large transit timing variations ('TTVs') of up to ±3 percent of any given trial period, indicates no significant periodicities. See Section 3 for details.](https://academic.oup.com/mnras/article-abstract/488/2/2455/5525096)
Table 4. BLS sensitivity tests$^a$.

| Injected transits | Period (d) | S/N |
|-------------------|------------|-----|
| 10                | 8.6        | 18.0|
| 9                 | 9.4        | 17.3|
| 8                 | 10.9       | 13.9|
| 7                 | 12.4       | 12.6|
| 6                 | 14.0       | 10.9|
| 5                 | 16.9       | 9.2 |
| 4                 | 22.0       | 7.2 |
| 3                 | 30.2       | 4.0 |

Note. $^a$Detection sensitivity to small numbers of periodic transits in the presence of 28 non-periodic ones.

Figure 3. The depth (in parts per million) in the flux of the 28 transit-like dips plotted against their duration in hours. There is no apparent correlation between these two parameters which characterize the dips.

In Fig. 3, we show a plot of the dip depth versus their duration. As is fairly obvious, there is no correlation between the two parameters. All but two depths lie in the range 200 ± 80 ppm.

We did one formal test for random arrival times among the 28 observed dips in flux. For a Poisson distribution of arrival times, the ‘survival’ probability for having no event following a given event for a time $\Delta t$ is proportional to $\exp\left(-\frac{\Delta t}{\tau}\right)$, where $\tau$ is the mean time between arrivals (3.1 d in our case). Since there are relatively few events to work with, and the values of $\Delta t$ cover two orders of magnitude, we decided to use the logarithm of $\Delta t$'s. The distribution of arrival time differences is then given by

$$dN/dz = N \exp[z - \exp(z)],$$

where $z \equiv \ln(\frac{\Delta t}{\tau})$, and $N$ is the number of arrival time differences. The distribution of arrival times (in logarithmic bins) and a fit to equation (1) is shown in Fig. 4. The only free parameter is the normalization, and that was computed using Cash statistics (Cash 1979) for the small numbers in each bin. The fit is good, considering the limited statistics. In the end, with only 28 events it is difficult to prove formally that a data set actually has random arrival times. And, it is likely that the arrival times of a several planet system would pass this same test (i.e. a fit to equation 1).

Finally, we created an average profile for the 28 dips by stacking them (see Fig. 5). Each dip is shifted to a common centre, and is then stretched (or shrunk) to a common duration of 2 h, and expanded (or contracted) vertically to a common depth of 200 ppm. It is not surprising, therefore, that the average profile has a duration close to 2 h and a depth near 200 ppm. However, what is more significant is that the average profile is closer to a planet-transit shape than to being triangular, and that no statistically significant asymmetries, long egress tail, or pre- or post-transit features are seen.

Figure 4. Distribution of the sequential interarrival times, $\Delta t$, of the 28 dips, expressed as the natural log of $\Delta t/\tau$, where $\tau$ is the mean rate of transit-like events (3.09 d per dip). The red fitted curve is the ‘exponential’ distribution, expressed in units of per logarithmic interval, that would be expected for random arrival times. The narrowly spaced vertical blue lines are what would be expected for a periodic function with 3 per cent TTVs.

Figure 5. Average profile of the 28 detected dips in EPIC 249706694 (HD 139139). The dips are shifted to a common centre, are stretched to a common duration of 2 h, and expanded vertically to a common depth of 200 ppm. One unit on the X and Y axes corresponds to 2 h and 200 ppm, respectively. No significant asymmetries, long egress tail, or pre- or post-transit features are seen (e.g. as in the cases of the so-called ‘disintegrating planets’; van Lieshout & Rappaport 2018.)

4 TESTS OF THE DATA SET

It is possible for the instrument systematics of the Kepler detector system to mimic planetary transits. In particular, the ‘rolling band artefact’ can cause small, transit-shaped changes in flux (see the Kepler Instrument Handbook; Thompson et al. 2016), and channel ‘cross-talk’ can produce spurious signals. Additionally, transiting signals originating from background targets could pollute the target and mimic transits in EPIC 249706694. We have vetted EPIC 249706694 against these common types of instrument systematic and performed a light centroid test to check for background
Based on our analysis, the signal is unlikely to be due to rolling band artefacts for the following reasons. First, rolling band events are usually infrequent, and low S/N, unlike the transits in EPIC 249706694. Secondly, since rolling band is an additive effect, background pixels are proportionally more affected by rolling band, and the background pixels of EPIC 249706694 do not show any transiting signals. Thirdly, we note that no rolling bands have been detected on the particular channel containing our target star (Van Cleve & Caldwell 2016) so that there is little risk that the features we are observing are due to rolling bands. However, out of an abundance of caution, we have also analysed the light curves of the 10 nearest neighbours to EPIC 249706694 and find no evidence of transits in either their light curves or background. These neighbour stars are listed in Table 5.

EPIC 249706694 is a bright target, and in order to create a spurious background signal of 200 ppm that would mimic a transit would require a significant contribution from an additive background. We find that all 10 of the nearest neighbours of EPIC 249706694 would have exhibited the equivalent, additive background signal at greater than 8σ confidence. As we do not detect any transits in background apertures or neighbouring stars, we rule out an additive, rolling band type signal as the cause of these dips.

In addition to the Kepler magnitudes of the neighbouring stars, Table 5 also gives the limiting depth (2σ) that could have been detected for dips in flux lasting for 3 h, were they to host a multiplicative, contaminating signal. Seven of the stars are faint, (Kp between 14.9 and 18.6), and we would not have achieved the sensitivity to detect the dips in EPIC 249706694. For three of the neighbours with Kp of 10.7 to 12.9, we would have the sensitivity to detect 200 ppm dips lasting for 3 h; however we detect no such dips in their light curves. This further suggests that the signal is not an instrument systematic.

It is possible to cause a spurious signal from detector `cross talk’, in which bright targets on one channel cause excess flux to be measured on other channels in the same module, at the same row/column coordinates. However, cross talk is a small effect, and requires a bright target on a channel in the same module. To create a spurious signal of 500 counts to cause these transiting signals would require a bright target on adjacent neighbouring channels. In this case, as EPIC 249706694 is on channel 42, to create the transiting signal would require a star of 10th magnitude on channels 41 or contaminants, and we find no evidence that the transiting signals are caused by instrument systematics or background stars.

Figure 6. Composite image of EPIC 249706694 showing our target star as well as two other EPIC stars 211831631 and 246253930, each observed during a different campaign. The red lines denote the postage-stamp boundaries for each star during its respective observation. We have verified that during observations of EPIC 211831631 and 246253930 the pixels in the overlap region were behaving normally.

We have checked the Pixel Response Function (PRF) centroids for EPIC 249706694 during the dips, as compared to out-of-dip regions. PRF centroid shifts during transits would indicate that the transiting signal is coming from a neighbouring or background target. We are able to establish that the PRF centroid during the transit is consistent with the PRF centroid out of transit, suggesting that the transiting signal originates from EPIC 249706694. To strengthen this interpretation, given that this star is saturated and bleeding, we formulated two sets of difference images of the average out-of-transit image versus the average in-transit image, where the out-of-transit data were taken to be one transit-duration on either side of each transit. In one set, we fitted out the measured motion of the star from each pixel prior to constructing the difference images to reduce the impact of the pointing deviations. In both cases, we found that many of the difference images exhibited features associated with the transits at the ends of the saturated bleed segments. In other cases, the difference image showed that the transit-like feature was occurring in the core of the stellar image on either the left or the right side of the saturated, bleeding columns. These difference images indicate that the source of the dips is collocated with, or very close to, the target star.

As reported in Table 2, a neighbour star exists within 3.3 arcsec of EPIC 249706694. Since this star exists in the bleed column of EPIC 249706694, we are not able to discount it as the origin of the transiting signals.

We have also checked the performance of the CCD pixels which detected the transit-like signals in EPIC 249706694 during observations of other stars during previous K2 campaigns. Fig. 6 shows a composite image of EPICs 249706694, 211831631, and 246253930 (the latter two from C5 and C12, respectively). Based on the measurements of these pixels from other campaigns, we have no evidence that there is a detector anomaly causing these

| Star (EPIC) | Kp | Limiting 3-h dip (ppm) | Separation (arcmin) |
|------------|----|------------------------|---------------------|
| 249704360  | 12.9 | 94                     | 1.8                 |
| 249708809  | 17.9 | 1372                   | 2.2                 |
| 249705070  | 10.7 | 32                     | 2.8                 |
| 249702101  | 17.8 | 1100                   | 3.8                 |
| 249704543  | 16.9 | 1586                   | 4.4                 |
| 249710690  | 16.7 | 1401                   | 5.3                 |
| 249700088  | 12.5 | 76                     | 5.3                 |
| 249713491  | 14.9 | 281                    | 6.2                 |
| 249713323  | 18.6 | 2014                   | 6.2                 |
| 249713897  | 17.9 | 4179                   | 6.3                 |

Notes. aAll the K2 observed stars within 6 arcmin of EPIC 249706694. b2σ sensitivity limit. cAngular separation between EPIC 249706694 and the neighbour star.
transit-like signals. However, the overlap in pixels is only partial, and so we cannot formally exclude an anomaly in a group of pixels that do not overlap the nearby C5 and C12 targets (Fig. 6). Known detector problems include the ‘Sudden Pixel Sensitivity Drop Out’ (‘SPSD’); see e.g. Christiansen et al. (2013) which causes permanent degradation, and would not lead to the transit-like behaviour we observe. Although we cannot completely rule out a detector anomaly as the source of the signal, we have no evidence to suggest that the pixels used to detect EPIC 249706694 are faulty, and no other reported detector anomaly shows a similar morphology to the transits we observe.

Finally, we have also considered the possibility that Solar system objects (e.g. asteroids) passing through or near the photometric aperture could somehow cause spurious dips in flux. We used SkyBot (Berthier et al. 2006, 2016) and find that at the arrival times of the dips (Table 1) there are no known catalogued asteroids within 1 arcmin of the target star (EPIC 249706694), except for the 16th dip in which we register the inner main belt asteroid 2009 BU179 at a distance of 36 arcsec. Specifically, no known asteroid passes through the photometric aperture which fits within a box that extends only 20 arcsec from the target star. Furthermore, on Kepler and K2 the background flux was determined via a grid of background pixel targets to which a two-dimensional polynomial was fitted, and then evaluated at each pixel location (Twicken et al. 2010), so it is at best unlikely that asteroids could produce spurious transits unique to a particular target. And, in this latter regard, we note that while 28 such dips were recorded for EPIC 249706694, none was found in any of the 10 neighbouring stars (see above discussion).

5 GROUND-BASED SPECTRA

5.1 TRES spectra

We observed EPIC 249706694 (HD 139139) with the Tillinghast Reflector Echelle Spectrograph (TRES; Fúrubisz, Szemgorygi & Meibom 2008) on the 1.5 m Tillinghast Reflector at the Fred Lawrence Whipple Observatory on Mt. Hopkins, Arizona. TRES has a spectral range of 3900–9100 Ångstroms and a resolving power of R ∼ 44 000. The spectra were reduced and extracted as described in Buchhave et al. (2010).

We obtained two radial velocity observations on UT 2018 April 8 and April 10. The spectra had an average S/N per resolution element (SNRe) of ≥35 at the peak continuum near the Mg b triplet at 519 nm with exposure times averaging 240 s. A multi-order velocity analysis was performed by cross-correlating the two spectra, order by order. 20 orders were used, excluding low S/N orders in the blue part of the spectrum and some red orders with contamination by telluric lines. The difference in the relative velocity was about a few m s\(^{-1}\), with an uncertainty of 36 m s\(^{-1}\), based on the rms scatter of the differences from the individual orders.

The two radial velocities that we obtained with TRES are listed in Table 6. Absolute radial velocities were derived for the two TRES spectra by cross-correlation against the best-matching synthetic template from the CFA library of calculated spectra using just the Mg b echelle order. The two spectra, taken two nights apart, yield the same RV to within 4 ± 100 m s\(^{-1}\), which is already quite constraining.

The Stellar Parameter Classification (SPC; Buchhave et al. 2010) tool was used to derive the stellar parameters of the target star. SPC cross correlates the observed spectra against a library of synthetic spectra based on Kurucz model atmospheres (Kurucz 1992). We calculated the average of the inferred parameters for the two spectra, and they are fully consistent with those given in the EPIC input catalogue (Huber et al. 2016). The stellar parameters for star A are summarized in Table 7.

5.2 McDonald spectra

We observed EPIC 249706694 with the high-resolution Tull spectrograph (Tull et al. 1995) on the 2.7 m telescope at McDonald Observatory in Ft. Davis, TX. We observed the target on 2018 April 6 using a 1.2 arcsec wide slit, yielding a resolving power of 60 000 over the optical band. We obtained three individual spectra back to back with 20-min exposures to aid in cosmic-ray rejection, which we combined in post-processing to yield a single higher S/N spectrum. We bracketed each set of three exposures with a calibration exposure of a ThAr arc lamp to determine the spectrograph’s wavelength solution. We extracted the spectra from the raw images and determined wavelength solutions using standard IRAF routines, and we measured the star’s absolute radial velocity using the Kea software package (Endl & Cochran 2016).

The McDonald spectrum of star A, taken 2 d earlier than the first TRES spectrum, agrees with the two TRES spectra to within 110 ± 160 m s\(^{-1}\). The stellar parameters for star A independently inferred from the McDonald spectrum using Kea are summarized in Table 7.

We also acquired three spectra of star B with the Tull spectrograph. It was difficult, however, to completely avoid the light
from the brighter nearby star A. None the less, we are able to extract RVs of 16.47 ± 0.48, 15.92 ± 0.35, and 16.06 ± 0.46 km s⁻¹ (see Table 6). These values are quite consistent, within the uncertainties, with the RV for star A and show no evidence for variability, indicating that star B is likely not a short-period binary.

5.3 Gaia spectra

The Gaia data base (Lindegren et al. 2018) lists an RV value for star A of 16.36 ± 0.40 km s⁻¹, which agrees with the two TRES spectra to within 200 ± 400 m s⁻¹. It also turns out that Gaia was able to get a rough radial velocity for star B which differs from that of star A by only 1 ± 4 km s⁻¹. We note, however, that the reported Gaia RV values are the median RVs for seven and six visits for star A and star B, respectively, over a 2-yr interval. The assigned uncertainty is roughly the standard error of the set of measurements. Thus, the upper limit to the K velocity of the binary motion for any orbital period that we might consider (e.g. P_orb < 1 yr) is less than 1 km s⁻¹ for star A and 10 km s⁻¹ for star B, based on the Gaia data alone.

5.4 RAVE spectra

The 5th data release of the RAAdial Velocity Experiment (‘RAVE’; Kunder et al. 2017) lists two RV measurements for star A taken in 2003 and 2008. They report heliocentric velocities of 19.7 ± 4.6 and 17.7 ± 1.1 km s⁻¹ for the two measurements, respectively. These are also listed in Table 6. Though not highly precise, these also indicate long-term stability in the radial velocity of star A.

5.5 Summary of spectral studies

From all the spectral observations discussed above, we draw two conclusions: (i) stars A and B likely form a bound binary (projected separation ~400 au) inferred from the similar RVs and distances, comparable proper motions, and close proximity on the sky; and (ii) neither star is itself in a close binary system because of lack of large RV variations.

We can further quantify the latter conclusion for star A. We utilized all the RV values listed in Table 6 to carry out a Fisher-matrix analysis (D. Wittman6) of the (2σ) upper limits to circular orbits with nearly 10⁶ trial periods between 0.5 d and 30 yr. We then adopted a mass for star A of 1.0 M⊙ and considered inclination angles from 30° to 90° in order to set corresponding limits on the mass of any unseen close companion to star A. The result is that stars and even brown dwarfs are ruled out for P_orb ≲ 10 d, low-mass stellar companions are disallowed for P_orb ≲ 100 d, and comparable mass stars are ruled out for periods below a few decades.

6 POSSIBLE EXPLANATIONS FOR THE DIPS

6.1 Planet transits in a multoplanet system

We have shown that no more than subsets of four dips could be caused by periodic transits. Let’s assume for the sake of argument that the host star has a chain of planets in 3:2 mean motion resonances. If the first has a period of 20 d, yielding four visible transits during C15, then the next planet could yield at most three transits, and so forth. Adding these up, we would see at most 13 transits before longer period planets began yielding, at best, a single improbable transit. It is farfetched to imagine that each of the remaining 15 transits could be accounted for by different individual planets, each more distant than the previous one, and still yielding transits all within an 87-d observation interval.

6.2 Planets orbiting both stars A and B

This hypothesis does not help much over that described in Section 6.1. It is also farfetched to imagine that each of star A and star B has a subset of the transits (e.g. 14 each), and only small subsets of each of these would be periodic – yet all of the transit depths are similar. This, in spite of the fact that the putative planets around star A would all have to be Earth size while all around star B would be Jovian size.

6.3 Few-planet system with huge TTVs

We also considered a case where the transits have sufficiently large TTVs so as not to be detectable in BLS searches. In that case, they should still be detectable in the IMT transform that we computed. In fact, we allowed TTVs as large as plus and minus 10 per cent of a period, and the IMT still did not find any underlying periodicity. We cannot envision a scenario where huge TTVs, e.g. causing a pattern such as seen in Fig. 1, can exist, and where the system remains dynamically stable.

6.4 Disintegrating planet with only rare transits appearing

Another possibility is an ultrashort period planet (‘USP’) where the transits are caused by dusty effluents rather than a hard body occultation (e.g. Rappaport et al. 2012; van Lieshout & Rappaport 2018). In that case, perhaps only a few of the transits would appear. Since the shortest intervals between ‘transits’ in the observed data set are ~5 h, this would have to be the maximum allowed period of the putative USP. In that case, there would have been more than 400 transits during the C15 observations. If only 28 of them are observed, then one would need to explain why the dust emissions were sufficiently active in only 7 per cent of the events to reveal a transit, but nearly completely absent in the remaining 93 per cent of cases. In any event, we have done simulations to show that a random selection of 10–15 per cent of several hundred periodic events is readily detectable in a BLS transform, but 7 per cent is getting close to undetectable. Finally, in this regard, we note that the average dip profile (see Fig. 5) shows no evidence for the type of asymmetries that might be expected from disintegrating planets.

6.5 Debris disc of dust-emitting asteroids

An initially attractive scenario is a debris disc hosting many dust-emitting asteroids or planetesimals over a substantial range of radial distances. Each transit-like event might be attributed to a single asteroid with dusty effluents transiting the host star. We note that the postulated dust-emitting asteroids in WD 1145 + 017 (Vanderburg et al. 2015) could produce 200 ppm transits when crossing a Sun-like star. The difficulties with such a scenario include: (i) the fact that almost all the transit-like events have rather similar depths, and (ii) it is doubtful that asteroids in a wide range of orbits would all be at just the right equilibrium temperatures to produce copious dust emission. These both point to the problem with numerous
independent bodies all having just the right properties to emit very nearly the same amounts of dust.

### 6.6 Eccentric planet orbit about one star in a close binary – S-type

One interesting possibility is that of a planet in an eccentric orbit about one star in a fairly close binary system, e.g. with a binary period of 2–5 d. The eccentric planetary orbit would require a very short orbital period of \( \lesssim 5–16 \) h to remain well inside the Roche lobe of the host star. The rapid precession of such an eccentric planetary orbit would allow for both different durations of the transits and a possible erratic set of arrival times. The putative binary system would have to be sufficiently tilted with respect to our line of sight so as to avoid producing binary eclipses – which are not observed. We have carried out a number of such dynamical simulations and found that we could never succeed in masking the underlying periodicity sufficiently to avoid detection in either a BLS or IMT transform. Furthermore, the radial velocity information we have (Table 6) also argues against a short-period binary, as does the fact that there is no sign of a composite spectrum in the TRES spectral observations.

### 6.7 P-type circumbinary planet

Another scenario we investigated is a P-type circumbinary planet. Transits across one or both stars and the stellar motions around the centre of mass may produce large variations in the transit timing and duration. This would be an extreme case of transiting circumbinary planets found by Kepler (e.g. Doyle et al. 2011; Welsh et al. 2012). To test this scenario, we performed the following simple analysis. We initialize the planetary orbit such that it eclipses one of the binary stars when the first transit event occurs at BKJD = 3161.6826. We ignored dynamical interactions in this simple model, i.e. we assumed that the binary and planetary motions follow independent Keplerian orbits with period ratios outside the stability limit (Holman & Wiegert 1999). To sample the orbital parameters, we define a likelihood function that is the sum of booleans of all individual K2 data points: (a) true if the time is within one of the transiting events (Table 1) and the planet is in front of one of the stars from the observer’s view in our model; (b) false if a transit is observed but the planet is not eclipsing in our model, or the planet eclipses in our model but no transit was observed. We minimized the corresponding likelihood function with a Levenberg–Marquardt algorithm to find the best solutions. However, after initializing the model with various starting parameters (\( P_{\text{orb}} \) typically near 2 d and \( P_{\text{plan}} \) near 8–10 d), we could not find a convincing solution that correctly predicted more than 5 of the events in Table 1. Given the crudeness of our model, the P-type circumbinary scenario is perhaps more decisively ruled out by the absence of RV variations in the host star (see Section 5).

### 6.8 Dipper-like activity

Some ‘dipper’ stars (Cody et al. 2014; Ansdell et al. 2016) exhibit quasi random dips that often range from a few to 50 per cent of the flux. If such dipper activity happens to take place in association with the fainter star B, then those dips would be diluted by about a factor of 15 due to the presence of star A which is in the same photometric aperture. To check the plausibility of this scenario, we have taken the light curves for the 10 dipper stars reported in Ansdell et al. (2016) (see their fig. 2), and have mocked up what one would have seen if the variability in them had been diluted by a factor of 15. In no case does the resultant ‘diluted’ version of the light curves closely resemble the light curve of the Random Transitor. In particular, at least one of the following is likely to apply: (i) there is significant activity in the out-of- ‘dip’ regions; or (ii) the host star has excess WISE band 3 and 4 emissions; or (iii) there is an underlying periodicity present comparable with the dip recurrence times. The fact that these are all different from the Random Transitor, and star A in particular, does not prove that ‘dipper’ phenomena are not involved in the transit-like events we see. For one example, star B might have a small WISE band 3 and 4 excess that we are unable to detect because of the dilution of star A. Finally, we note that since there is currently no quantitative model for explaining dipper stars, invoking dipper behaviour to explain the Random Transitor does not provide any real insight to the mystery at hand.

### 6.9 Short-lived star spot

Finally, we considered cases where the dips in flux are caused by a short-lived change in the intrinsic flux from the host star. One idea is that the host star has short-lived spots that are larger than the Earth in size on star A, or larger than Jupiter on star B, but appear suddenly (i.e. within a half hour), last for a few hours, and then disappear equally suddenly. We do not believe that this violates any physical laws. Furthermore, the existence of such spot activity on distant stars studied with Kepler, K2, and TESS, might have gone either unnoticed or unreported by Citizen Scientists or by automated searches for single-transit systems. This is especially true if such dips are much less frequent than once per month, are aperiodic, are at the 200 ppm level, and last for only a couple of data cadences. In the latter case, they might be considered too short for the implied longer periods associated with single-transit events.

### 7 SUMMARY AND CONCLUSIONS

K2 has provided us with a large number of impressive results. Several mysterious objects, however, linger even after the end of the mission. In this work, we have presented a K2 target, EPIC 249706694 (HD 139139), that exhibits 28 transit-like events that appear to occur randomly. The depths, durations, and times of these dips are shown in Figs 1 and 3 and summarized in Table 1. Most of the events have depths of 200 ± 80 ppm, durations of 1–7 h, and a mean separation time of 3 d.

The target star EPIC 249706694 (star A) is a relatively unevolved star of spectral type G. It has a rotation period of \( \approx 14.5 \) d, which implies a gyrochronological age of \( 1.5 \pm 0.4 \) Gyr (Barnes 2007). However, in the same photometric aperture is a fainter neighbouring star (3.3 arcsec away) that is 2.7 \( G_p \) magnitudes fainter with \( T_{\text{eff}} \approx 4400 \) K. There is insufficient Gaia kinematic data for this star (B) to know for sure if it is bound to star A, but its similar RV value, distance, and estimated close relative proper motion (see Section 2.3) imply that it is likely a physically bound companion to star A. In any case, we do not know with certainty which star hosts the transit-like events. If the dips are transits across star A then the inferred planet sizes would be \( \approx 1.5 \) \( R_p \), whereas if they are on star B, then the planetary radii would be more like Jupiter’s size.

Neither star A nor star B appears to exhibit large RV variations, with limits of \( \lesssim 1 \) km s\(^{-1}\) (see Section 5 and Table 6). The constraints that the RV measurements place on possible orbital periods for binary companions are discussed at the end of Section 5. For star

---

\( \text{https://keplerscience.arc.nasa.gov/scicon-2019/} \)
A. we find that stellar- and even brown-dwarfs-companions are ruled out for \( P_{\text{orb}} \lesssim 10 \) d, and low-mass stellar companions are disallowed for \( P_{\text{orb}} \lesssim 100 \) d. For star B the results are somewhat less constraining.

We have done some exhaustive tests to ensure that the observed shallow dips in the flux of EPIC 249706694 are of astrophysical origin (see Section 4). These include ruling out rolling-band artefacts, centroid motions during the dips, cross-talk, or bad CCD pixels. One can never be entirely sure that such transit-like events are astrophysical, but we have done our best due diligence in this regard.

We have briefly considered a number of scenarios for what might cause the randomly occurring transit-like events. These include actual planet transits due to multiple or dust emitting planets, dust-emitting asteroids, one or more planets with huge TTVs, S- and P-type transits in binary systems, ‘dipper’ activity, and short-lived spot activity. We find that none of these, though intriguing, is entirely satisfactory. These are reviewed in Section 6.

The purpose of this paper is largely to bring this enigmatic object to the attention of the larger astrophysics community in the hope that (i) some time on larger telescopes, or ones with high photometric precision, might be devoted to its study, and (ii) some new ideas might be generated to explain the mysterious dips in flux. With regard to observational advances, two specific things could be done. First, acquire more and higher quality spectra of stars A and B to improve the constraints on RV variability. Secondly, if time is available on a telescope with 1 mmag precision photometry in the exposure time exceeding a few days. There would then be a chance of directly measuring the transit-like events on star B – if that is indeed their origin.

ACKNOWLEDGEMENTS

AV’s and KM’s work was supported in part under a contract with the California Institute of Technology (Caltech)/Jet Propulsion Laboratory (JPL) funded by National Aeronautics and Space Administration (NASA) through the NASA Exoplanet Science Institute. We thank Allan R. Schmitt and Troy Winarski for making their light curve examining software available on a telescope with 1 mmag precision photometry in the presence of a brighter star 3.3 arcsec away, monitor star B for a total exposure time exceeding a few days. There would then be a chance of directly measuring the transit-like events on star B – if that is indeed their origin.

REFERENCES

Ansdell M. et al., 2016, ApJ, 816, 69
Bailer-Jones C. A. L., Rybizki J., Fouesneau M., Mantelet G., Andrae R., 2018, AJ, 156, 58
Barensten G., Cardoso J. V. d. M., 2018, Astrophysics Source Code Library, record ascl:soft03005B
Barnes S. A., 2007, ApJ, 669, 1167
Becker J. C., Vanderburg A., Adams F. C., Rappaport S., Schwengeler H. M., 2015, ApJ, 812, L18
Berthier J., Vachier F., Thuillot W., Fernique P., Ochsenbein F., Genova F., Lainey V., Arlot J.-E., 2006, ASP Conf. Ser. Vol. 351, Astronomical Data Analysis Software and Systems XV. Astron. Soc. Pac., San Francisco, p. 367
Berthier J., Carry B., Vachier F., Eggl S., Santerne A., 2016, MNRAS, 458, 3394
Borkovits T. et al., 2018, MNRAS, 478, 5135
Borkovits T. et al., 2019, MNRAS, 483, 1934
Buchhave L. A. et al., 2010, ApJ, 720, 1118
Cash W., 1979, ApJ, 228, 939
Chaplin W. J. et al., 2015, PASP, 127, 1038
Christiansen J. L. et al., 2013, ApJS, 207, 35
Cody A. M. et al., 2014, ApJ, 147, 82
Curi R. M. et al., 2013, wise.rept, 1C
Dattilo A. et al., 2019, AJ, 157, 169
Dimitriadi G. et al., 2019, ApJ, 870, 1
Doyle L. R. et al., 2011, Science, 333, 1602
Endl M., Cochran W. D., 2016, PASP, 128, 94502
Flewelling H. A. et al., 2006, preprint (arXiv:1612.05243)
Furész G., Szentygyorgyi A. H., Meibom S., 2008, in Santos N. C., Pasquini L., Correia A. C. M., Romaniello M., eds, Precision Spectroscopy in Astrophysics, Vol. 287, ESO Astrophysics Symposia book series (ESO). Springer, p. 290
Gaidos E. et al., 2014, MNRAS, 443, 2561
Garnavich P., 2019, Better Understanding Supernovae from Kepler/K2 Observations, talk presented at the Fifth Kepler & K2: Kepler Science Center Website
Gay B. L., Rappaport S., Kaye T. G., Alonso R., Hambschs F.-J., 2017, MNRAS, 465, 3267
Grindlay J. T., Simcoe R., Laycock S., 2009, in Osborn W., Robbins L., eds, ASP Conf. Ser. Vol. 410, Preserving Astronomy’s Photographic Legacy: Current State and the Future of North American Astronomical Plates. Astron. Soc. Pac., San Francisco, p. 101
Holman M. J., Wiepert P. A., 1999, AJ, 117, 621
Howell S. B. et al., 2014, PASP, 126, 398
Huber D. et al., 2016, ApJ, 224, 2
Jayasinghe T. et al., 2018, MNRAS, 477, 3145
Jenkins J. M. et al., 2018, MNRAS, 485, 2681
Kipping D. M., Schmitt A. R., Huang X., Torres G., Nesvorný D., Buchhave L. A., Hartman J., Bakos G. Á., 2015, ApJ, 813, 14
Koch D. G. et al., 2010, ApJ, 713, L79
Kovács G., Zucker S., Mazeh T., 2002, A&A, 391, 369
Kunder A. et al., 2017, AJ, 153, 75
Kurucz R. L., 1992, in Barbuy B., eds, ASP Conf. Ser. Vol. 149, The Stellar Populations of Galaxies. Angra dos reis, Brazil, p. 225
Leiner E., Mathieu R. D., Vanderburg A., Gosnell N. M., Smith J. C., 2019, preprint (arXiv:1904.02169)
Lindegren L. et al., 2018, A&A, 616, 25
Lugr R., Agol E., Kruse E., Barnes R., Becker A., Foreman-Mackey D., Deming D., 2016, AJ, 152, 100L
Lund M. N. et al., 2016a, PASP, 128, 12404
Lund M. N. et al., 2016b, MNRAS, 463, 2600
Mayo A. W. et al., 2018, AJ, 155, 136
Molnár L. et al., 2018, ApJS, 234, 37
Rappaport S. et al., 2017, MNRAS, 467, 2160
Rappaport S. et al., 2019, MNRAS, 485, 2681
Rappaport S. A. et al., 2012, ApJ, 752, 1
Rest A. et al., 2018, Nature Astron., 2, 307
Ricker G., Winn J. N., Vanderspek R., 2014, in Oschmann Jr. J. M., Clampin M., Fazio G. G., MacEwen H. A., eds, Proc. Conf. Ser. Vol. 9143, Space Telescopes and Instrumentation 2014: Optical, Infrared, and Millimeter Wave. SPIE, Bellingham, p. 20
Rodriguez J. E. et al., 2018, AJ, 156
Sanitas-Ojeda R. et al., 2015, ApJ, 812, 112
Skrutskie M. F. et al., 2006, AJ, 131, 1163.
Stello D. et al., 2017, ApJ, 835, 83
Thompson S. E., Fraquelli D., van Cleve J. E., Caldwell D. A., 2016, Kepler Archive Manual (KDMC-10008-006)
Toymassian G., Szkody P., Yarza R., Kennedy M., 2018, ApJ, 863, 47
Tull R. G., MacQueen P. J., Sneden C., Lambert D. L., 1995, PASP, 107, 251
Twicken J. D., Clarke B. D., Bryson S. T., 2010, Radziwill N. M., Bridger A., Proc. Conf. Ser. Vol. 7740, Software and Cyberinfrastructure for Astronomy. SPIE, Bellingham, 774023
Van Cleve J. E., Caldwell D. A., 2016, in Haas M. R., Howell S. B., eds, Kepler Science Document, KSCI-19033-002
van Lieshout R., Rappaport S., 2018, in Deeg H. J., Belmonte J. A., eds, Handbook of Exoplanets. Springer-Verlag, Berlin
Vanderburg A., Johnson J. A., 2014, PASP, 126, 948
Vanderburg A. et al., 2015, Nature, 526, 546
Vanderburg A. et al., 2016, ApJ, 827, 10
Welsh W. F. et al., 2012, Nature, 481, 475
Zhou G. et al., 2018, ApJ, 854, 109
Zhu W. et al., 2017, ApJ, 849, 31

This paper has been typeset from a \TeX/S\LaTeX file prepared by the author.