ACCRETION IN THE EARLY OUTER SOLAR SYSTEM

SCOTT J. KENYON
Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138; skenyon@cfa.harvard.edu

AND

JANE X. LUU
Leiden Observatory, P.O. Box 9513, 2300 RA Leiden, Netherlands; luu@strw.leidenuniv.nl

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ABSTRACT

We describe calculations of the evolution of an ensemble of small planetesimals in the outer solar system. In a solar nebula with a mass of several times the minimum mass solar nebula, objects with radii of 100–1000 km can form on timescales of 10–100 Myr. Model luminosity functions derived from these calculations agree with current observations of bodies beyond the orbit of Neptune (Kuiper Belt objects). New surveys with current and planned instruments can place better constraints on the mass and dynamics of the solar nebula by measuring the luminosity function at red magnitudes, $m_R \geq 28$.

Subject headings: solar system: formation — Kuiper Belt, Oort Cloud

1. INTRODUCTION

Several remarkable discoveries have renewed interest in solar system formation. Recent surveys have detected many small icy bodies beyond the orbit of Neptune (e.g., Jewitt & Luu 1993; Williams et al. 1995; Jewitt, Luu, & Chen 1996; Luu et al. 1997; Gladman & Kavelaars 1997; Chiang & Brown 1999). Assuming a geometric albedo of 4%, these Kuiper Belt objects (KBOs) have radii of 50–400 km; the derived size distribution implies a significant population of smaller objects. With semimajor axes of 40–50 AU and orbital inclinations of $0^\circ$–$30^\circ$, the orbits of known KBOs suggest an annulus of planetesimals formed in situ and left over from the planetary formation epoch (Holman & Wisdom 1993). The presumed structure of this annulus resembles the dusty disks recently discovered around several nearby stars (Smith & Terrile 1984; Aumann et al. 1984; Jayawardhana et al. 1998; Koerner et al. 1998; Greaves et al. 1998). Planets similar to those in our solar system have not been detected in any of these disks, but direct images and radial velocity measurements of nearby stars already imply the existence of more than one dozen extrasolar planets of several Jupiter masses (Latham et al. 1989; Marcy & Butler 1996; Cochran et al. 1997; Noyes et al. 1997; DelFosse et al. 1998; for a review see Marcy 1999).

These discoveries challenge planet formation theories. Most theories presume that planets grow by accretion of small planetesimals in a gaseous circumstellar disk (Safronov 1969; Goldreich & Ward 1973; see also Lissauer & Stewart 1993, Boss 1997, and references therein). Hitherto, numerical studies have focused on the formation of the prototypical terrestrial and gas giant planets, Earth and Jupiter (Greenberg et al. 1978, 1984; Nakagawa, Hayashi, & Nakazawa 1983; Ohtsuki, Nakagawa, & Nakazawa 1988; Wetherill & Stewart 1989, 1993; Barge & Pellat 1990; Ruden & Pollack 1991; Pollack et al. 1996; Weidenschilling et al. 1997). If the initial disk mass is comparable to the minimum mass solar nebula, these calculations often have difficulty producing objects similar to the known terrestrial or gas giant planets during the estimated disk lifetime of $\sim 10$–30 Myr (Pollack et al. 1996; Weidenschilling et al. 1997). This problem is exacerbated in the outer solar system, where numerical calculations yield formation times exceeding 100 Myr for 500–1000 km radius KBOs (Stern 1995, 1996; Stern & Colwell 1997a, 1997b). KBOs must form on shorter timescales in parallel with Neptune. Otherwise, Neptune’s gravity increases the velocities of nearby planetesimals, including those in the inner Kuiper Belt, on timescales of $\sim 10$ Myr (Malhotra 1996). This process prevents the growth of KBOs with radii exceeding 100–200 km, because large velocities hinder agglomeration.

We recently began to consider KBO formation in the outer solar system using an evolution code that follows planetesimal growth in the annulus of a circumstellar disk. Initial results indicate that KBOs can form at 30–50 AU on timescales of 10–100 Myr in disks with 1–3 times the minimum mass solar nebula when collisional disruption of planetesimals is unimportant (Kenyon & Luu 1998, hereafter KL98). Further calculations with an algorithm that includes disruptive processes lead to similarly short timescales for a wide range of initial conditions (Kenyon & Luu 1999; hereafter KL99). Here, we briefly summarize these new results, compare the theoretical model with current observations, and make predictions for comparison with future observations of KBOs.

$^1$ The minimum mass solar nebula has a surface density $\Sigma = \left(\frac{\Sigma_0}{R_{\oplus}}\right)^{3/2}$, where $\Sigma_0$ is the surface density of solid material at $R_{\oplus} = 1$ AU. We adopt $\Sigma_0 = 45 \text{ g cm}^{-2}$ (Hayashi 1981; see also Weidenschilling 1977; Bailey 1994). This definition yields a total mass of solids $M_0 \approx 100 M_\oplus$ inside the orbit of Neptune ($R < 30$ AU) and $M_0 \approx 10 M_\oplus$ in the inner part of the Kuiper Belt ($R = 32$–38 AU), where $1 M_\oplus = 6 \times 10^{27}$ g is the mass of the Earth.
TABLE 1

| Parameter                   | Symbol | Value |
|-----------------------------|--------|-------|
| Width of annulus (AU)       | $\delta a$ | 6     |
| Initial velocity (m s$^{-1}$) | $V_0$  | 0.45–45 |
| Particle mass density (g cm$^{-3}$) | $\rho_0$ | 1.5   |
| Relative gas velocity (m s$^{-1}$) | $\eta$ | 30    |
| Time step (yr)              | $\delta t$ | 1–250 |
| Number of mass bins         | $N$    | 64–256 |
| Mass spacing of bins        | $\delta$ | 1.25–2.0 |
| Minimum velocity for cratering (cm s$^{-1}$) | $V_f$ | 1     |
| Impact strength (ergs g$^{-1}$) | $S_0$  | $2 \times 10^6$ |
| Crushing energy (ergs g$^{-1}$) | $Q_c$ | $5 \times 10^4$ |
| Fraction of KE in ejecta    | $f_{ke}$ | 0.05  |
| Coefficient of restitution  | $c_1$  | $10^{-2}$ |
| Coefficient of restitution  | $c_2$  | $10^{-3}$ |

2. MODEL

Our accretion code is based on the particle-in-a-box method, where planetesimals are a statistical ensemble of bodies with a distribution of horizontal and vertical velocities about Keplerian orbits (Safronov 1969). We perform calculations for a single annulus of width $\Delta a$ centered at a heliocentric distance $a$. We approximate the continuous distribution of particle masses with $i$ discrete batches having particle populations $n_i(t)$ and total masses $M_i(t)$. The horizontal and vertical velocity dispersions are $h_i(t)$ and $v_i(t)$ (Wetherill & Stewart 1993). The average mass of a batch, $m_i(t) = M_i(t)/n_i(t)$, changes with time as collisions add and remove bodies from the batch. This procedure conserves mass and provides a statistical method to follow the growth of $\gtrsim 10^{20}$ small planetesimals into a few planets. Detailed $n$-body calculations confirm the basic features of particle-in-a-box calculations for the early stages of planet growth described here (Ida & Makino 1992; Kokubo & Ida 1996).

To evolve the initial size distribution in time, we calculate collision rates for the coagulation equation, determine the outcome of each collision, and compute velocity changes due to collisions and long-range gravitational interactions (see KL99). Each two-body collision can produce (1) merger into a single body with no escaping debris (very low impact velocity), (2) merger into a single body with escaping debris ("cratering"); low impact velocity), (3) rebound with or without cratering (modest impact velocity), or (4) catastrophic disruption into numerous smaller bodies (high impact velocity). The collision outcomes depend on the ratio of the impact energy $Q_f$ to the disruption energy $Q_d$ of two colliding planetesimals (Greenberg et al. 1978; Wetherill & Stewart 1993; Davis, Ryan, & Farinella 1994). Collisions with $Q_f > Q_d$ disrupt planetesimals into many small fragments. Collisions with $Q_f < Q_d$ yield a merged planetesimal and some small fragments if the collision velocity $V_f$ exceeds the minimum velocity for cratering $V_c$. Collisions with $Q_f < Q_d$ and $V_c < V_f$ yield a merged planetesimal with no cratering debris. We use an energy-scaling formalism to compute $Q_d$ as the sum of the intrinsic tensile strength $S_0$ and the gravitational binding energy (Davis et al. 1985, 1994). The intrinsic strength is the dominant component of $Q_d$ for bodies with $r_i \lesssim 1$ km; gravitational binding dominates $S_0$ for larger bodies.

For each collision, a velocity evolution algorithm distributes the kinetic energy among the resulting bodies and then accounts for collisional damping, kinetic energy transfer during elastic collisions ("dynamical friction"), angular momentum transfer during elastic collisions ("viscous stirring"), and gas drag (Hornung, Pellat, & Barge 1985; see also Wetherill & Stewart 1993; KL98). Dynamical friction tries to enforce equipartition of kinetic energy between mass batches; viscous stirring increases the velocities of all bodies. Gas drag removes objects from the annulus and reduces the velocities of small objects that are well coupled to gas in the disk.

We tested the code against analytical solutions and published numerical results (KL98; KL99). We reproduced previous calculations for accretion at 1 AU (Wetherill & Stewart 1993) and collisional disruption of preexisting large KBOs at 40 AU (Davis & Farinella 1997). Our calculations match analytical solutions well when the mass spacing between successive batches, $\delta = m_i + 1/m_i = 1.1–1.4$. Numerical solutions lag the analytic results by $\sim 10\%$ when $\delta = 1.4–2$. The timescales to produce objects of a given size increases with $\delta$, because poorer resolution prevents growth of large objects (see KL98).

Table 1 lists basic input parameters. The input cumulative size distribution $N_C$ has the form $N_C \propto r_i^n$, with initial radii $r_i = 1–80$ m. The total mass in the annulus is $M_0$; $M_0 \approx 10 M_\oplus$ for a minimum mass solar nebula. All batches start with the same initial velocity. We tested a range of initial velocities corresponding to initial eccentricities of $e_i = 10^{-4}$ to $10^{-2}$, as is expected for planetesimals in the early solar nebula (Malhotra 1995). The adopted mass density, $\rho_0 = 1.5 \text{ g cm}^{-3}$, is appropriate for icy bodies with a small rocky component. The fragmentation parameters, $V_f$, $Q_d$, $Q_c$, $h_i$, $c_1$, and $c_2$, are adopted from earlier work. KL99 describe these parameters in more detail.

To provide observational constraints on the models, we note that the known Kuiper Belt population contains at least one body with a radius of $\sim 1000$ km (Pluto) and $\sim 10^3$ KBOs with radii $r_i \gtrsim 50$ km between 30–50 AU. The cumulative size distribution of known KBOs can be fitted with $N_C \propto r_i^2$, with $q_{obs} = 3 \pm 0.5$ (Jewitt, Luu, & Trujillo 1998; see also Chiang & Brown 1999). Successful models should reproduce these observations on timescales comparable to (1) the estimated lifetimes of the solar nebula and gaseous disks surrounding nearby young stars, $\lesssim 10^7$ yr (Russell et al. 1996; Hartmann et al. 1998) and (2) the formation timescale for Neptune, $\lesssim 10^5$ yr (Lissauer et al. 1996).

3. NUMERICAL RESULTS

We separate the growth of KBOs into three regimes. Early in the evolution, frequent collisions damp the velocity dispersions of small bodies. These bodies slowly grow into 1 km objects on a timescale that is approximated by $\tau_{km} \approx 8 \text{ Myr} (M_0/10 M_\oplus)^{-1} (e_i/10^{-3})^{-0.65}$. This linear growth phase ends when the gravitational range of the largest objects exceeds their geometric cross section. This "gravitational focusing" enhances the collision rate by factors of 10–1000. The largest objects then begin a period of "runaway growth," when their radii grow from $\sim 1$ km to $\gtrsim 100$ km in several Myr. During this phase, dynamical friction and viscous stirring increase the velocity dispersions of the smallest bodies from $\sim 1$ m s$^{-1}$ up to $\sim 40$ m s$^{-1}$. This velocity evolution reduces gravitational focusing factors and ends runaway growth. The largest objects then grow slowly to $1000+ \text{ km}$ sizes on timescales that again depend on $M_0$. 

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and \( e_0 \). Column (5) in Table 2 lists timescales to form Pluto-sized objects \( \tau_P \) as a function of the input parameters \( M_o, \delta, e_0, \) and \( q_0 \).

Figure 1 shows cumulative size distributions for a model with \( M_o = 10 \, M_\oplus, q_0 = 3, \) and \( S_0 = 2 \times 10^6 \) ergs g\(^{-1}\). The shapes of these curves depend on two competing physical processes: (1) growth by mergers and (2) erosion by high-velocity collisions. In this example, collisions result in growth because the velocity dispersion is less than the catastrophic disruption threshold. However, the collision velocity exceeds the cratering threshold \( V_f \). Cratering adds debris to all low-mass batches. Gas drag removes material from low-mass batches \( (r_i \lesssim 10 \, \text{m}) \) but is ineffective at removing larger objects. The size distribution thus becomes shallower at small masses. At large masses, mergers produce a group of growing planetesimals with a steep size distribution. Once gravitational focusing becomes effective, the largest of these objects “run away” from the rest of the ensemble to produce a smooth power law with a maximum radius \( r_{\text{max}} \). As the evolution proceeds, \( r_{\text{max}} \) increases, but the slope of the smooth power law remains nearly constant.

The main features of these results depend little on the input parameters. All calculations produce two cumulative power-law size distributions connected by a transition region having an “excess” of planetesimals (the “bump” in the curves in Figure 1). The characteristic radius of this transition region increases from 0.3 km at \( e_0 = 10^{-4} \) to 3 km at \( e_0 = 10^{-2} \). If fitted with a power law of the form \( N_C \propto r_f^{-q_f} \) at small masses, the cumulative size distribution follows the predicted limit for collisional evolution, \( q_f = 2.5 \) (Dohnanyi 1969). We perform least-squares fits to obtain \( q_f \) at larger masses; column (6) of Table 2 lists derived values for \( q_f \) along with the 1 σ error. Column (7) lists the radius range for each fit. The results are surprisingly independent of the input parameters. We find the small range \( q_f = 2.75–3.25 \) for calculations with \( M_o = 100 \, M_\oplus, e_0 = 10^{-4–10^{-2}}, q_0 = 1.5–4.5, \) and \( S_0 = 10–3 \times 10^6 \) ergs g\(^{-1}\). This model result is consistent with the observed slope, \( q_{\text{obs}} = 3 \pm 0.5 \) (e.g., Jewitt et al. 1998).

### 4. Comparisons with Observations

As shown in Table 2, several calculations meet the success criteria defined in § 2. Annuli with \( M_o \gtrsim 10 \, M_\oplus \) produce Pluto-sized objects on short timescales, \( \tau_P \lesssim 50 \, \text{Myr} \) (for \( e_0 \lesssim 10^{-3} \)). Models with smaller initial masses or larger initial eccentricities form Plutos on longer timescales, \( \tau_P \gtrsim 50 \, \text{Myr} \). Plausible ranges of other input parameters, such as \( q_0, S_0, \) and \( f_{KB} \), yield \( \pm 20\% \) variations about these timescales. The results are insensitive to \( V_f \) and other collision parameters (KL99).

The crosses in Figure 1 compare our calculations directly with several observational constraints. The cross at \( r_i = 50 \, \text{km} \) indicates the number of KBOs with \( r_i \gtrsim 50 \, \text{km} \) estimated from recent ground-based surveys (Jewitt et al. 1998; see also Chiang & Brown 1999); the one at \( r_i = 10 \, \text{km} \) shows limits derived from a single, controversial measurement with Hubble Space Telescope (HST; Cochran et al. 1995, 1998; Brown, Kulkarni, & Liggett 1997). The third

![Figure 1](image-url)
cross plots limits at $r_i = 1$ km derived from theoretical attempts to explain the frequency of short-period comets from the Kuiper Belt (Davis & Farinella 1997; Duncan & Levison 1997; Levison & Duncan 1997). Our predictions agree with ground-based surveys at 50 km and theoretical limits at 1 km, but fall a factor of $\sim 10$ short of the HST measurement at 10 km.

To compare with observations in more detail, we predict the luminosity function (LF) of KBOs directly from the computed number distribution. We use a Monte Carlo calculation of objects selected randomly from the cumulative size distribution $N_C$. We assign each object a distance from the Sun $d_\odot$ and a random phase angle $\beta$ between the line of sight from the Earth to the object and the line of sight from the Sun to the object. This phase angle lies between $0^\circ$ and a maximum phase angle that is distance dependent. The distance of the object from the Earth is then $d_\oplus = d_\odot \cos \beta - [1 + d_\odot^2 \cos^2 \beta - 1]^{1/2}$. We derive the red magnitude of this object from a two-parameter magnitude relation for asteroids, $m_{R, KBO} = R_0 + 2.5 \log \left( t_i / t_2 \right) - 5 \log r_{KBO}$, where $R_0$ is the zero point of the magnitude scale, $r_{KBO}$ is the radius of the KBO, $t_1 = 2 d_\odot d_\oplus$, and $t_2 = \omega[(1 - g)\phi_1 + g\phi_2]$ (Bowell et al. 1989). In this last expression, $\omega$ is the albedo and $g$ is the slope parameter; $\phi_1$ and $\phi_2$ are phase functions that describe the visibility of the illuminated hemisphere of the object as a function of $\beta$. We adopt standard values, $\omega = 0.04$ and $g = 0.15$, appropriate for comet nuclei (Jewitt et al. 1998). The zero point $R_0$ is the apparent red magnitude of the Sun, $m_{R, \odot} = -27.11$, with a correction for the $V-R$ color of a KBO, $R_0 = m_{R, \odot} + \delta(V-R)_{KBO}$. Observations suggest that KBOs have colors that range from roughly $-0.1$ to $0.3$ mag redder than the Sun. We treat this uncertainty by allowing the color to vary randomly in this range.

The important parameters in the model LF are the distributions of input sizes (derived from the accretion calculations), distances, and orbital parameters. We assume KBOs are evenly distributed between “plutinos,” objects in 3:2 orbital resonance with Neptune having semimajor axes of $39.4 \pm 0.2$ AU, and “classical” KBOs with semimajor axes between 42–50 AU. The distance parameters are set by observations (Jewitt et al. 1998). This distance distribution is different from the 32–38 AU adopted for the coagulation calculations. Several tests show that accretion results at 42–50 AU are identical to those at 32–38 AU, except that the timescale to produce Pluto-sized objects is 50%–100% longer. To compute the model LF from the Monte Carlo magnitude distribution of KBOs, we scale the mass in the 32–38 AU annulus to match the mass in a 42–50 AU annulus, add in an equal number of plutinos, and divide by the sky area. The distribution of KBO orbital parameters is poorly known. We adopt circular orbits to derive magnitudes; the LF is insensitive to other choices. We assume orbital inclinations of $i = 0^\circ$–$5^\circ$ to compute the sky area, which is a compromise between the $i \approx 0^\circ$–$5^\circ$ of classical KBOs and the $i \approx 10^\circ$–$30^\circ$ of plutinos. The model LFs scale inversely with $\sin i$.

Figure 2 compares several models with the observed LF. The left panel shows models with $e_0 = 10^{-3}$ and different masses; the right panel shows models with the mass of a

**Fig. 2.—**Comparison of model luminosity functions of KBOs with observations. Data are as indicated in the legend of each panel. The open circle with the central dot indicates the position of Pluto for an adopted albedo of 4%; other observations are from Cochran et al. (1998; with HST), Irwin, Tremaine, & Zytkow (1995; 1995), Kowal (1989; K89), Luu & Jewitt (1988; LJ88), Gladman et al. (1998; G98), Luu & Jewitt (1998; LJ98), Jewitt et al. (1998; JLT98), and Chiang & Brown (1999; CB99). For clarity, error bars for each datum (typically a factor of 2–3) and the upper limit from Levison & Duncan (1990) are not shown. The lines plot luminosity functions for models with (left panel) $e_0 = 10^{-3}$ and $M_0 \approx 0.3$ (dot-dashed curve), 1.0 (solid curve), and 3.0 (dashed curve) times the minimum mass solar nebula and (right panel) a minimum mass solar nebula with $e_0 = 10^{-3}$ (dashed curve), $e_0 = 10^{-2}$ (solid curve), and $e_0 = 10^{-4}$ (dot-dashed curve). A minimum mass solar nebula has $M_0 \approx 12 M_\oplus$ within $R = 42–50$ AU. The pair of vertical solid lines indicates the planned magnitude range accessible to NGST.
minimum mass solar nebula and different $e_0$. Model LFs with the minimum mass and any $e_0$ agree with current observations. The good agreement of all models at $m_R \leq 20$, where the uncertainties are largest, depends on the assumed maximum radius in the model distribution. We picked 1000 km for convenience. Model LFs for $m_R \geq 20$ are independent of this choice.

To quantify the comparison between models and observations, we fit model LFs to $\log \mathcal{L}(m_R) = \alpha (m_R - m_0)$ over $20.5 \leq m_R \leq 26.5$. Table 3 lists the fitted $\alpha$ and $m_0$ as a function of the mass in classical KBOs (in units of the minimum mass solar nebula), $e_0$, the inner annulus boundary $R_{in}$, and the outer annulus boundary $R_{out}$. The small range in $\alpha$ for model LFs agrees with published values derived from observations. The $\alpha$ is independent of the relative numbers of plutinos and classical KBOs and of the distance distribution of classical KBOs. The observed zero point of the LF, $m_0 \approx 23.2-23.5$, favors models with masses comparable to the minimum mass solar nebula and any initial eccentricity. These data rule out models with $\leq 30\%$ of the minimum mass at the 3 $\sigma$ level. Smaller plutino fractions require larger masses: if plutinos are $10\%-25\%$ of the total KBO population, as indicated by recent observations (Jewitt et al., 1998), the needed mass is $2-4$ times the minimum mass.

There are two main uncertainties in comparing our model LFs with the data: the evolution of the KBO LF with time and the current orbital parameters of KBOs. The initial mass in KBOs was larger than implied by a direct comparison between the data and model LFs, because high-velocity collisions and dynamical encounters with Neptune have eroded the Kuiper Belt over time (Holman & Wisdom 1993; Davis & Farinella 1997; see also Levison & Duncan 1993; Duncan, Levison, & Budd 1995). Erosion from high-velocity collisions probably does not change the slope of the LF significantly. Massive KBOs with $r_i \gtrsim 50$ km ($m_R \lesssim 26-27$) are probably safe from collisional disruption (Davis & Farinella 1997). Disruption of smaller bodies depends on the unknown bulk properties and the poorly known orbital parameters of KBOs. These uncertainties are not important for comparisons of models and observations for $m_R \lesssim 26-27$ but can bias future comparisons.

Gladman et al. (1998) report $\alpha = 0.76^{+0.10}_{-0.11}$ and $m_0 = 23.4^{+0.20}_{-0.19}$ from a maximum likelihood analysis of previous surveys with magnitude limits, $20 \leq m_R \leq 28$. Surface densities in their Table 3 yield $\alpha \approx 0.6$ and $m_0 \approx 22.4$. Jewitt et al. (1998; see also Luu & Jewitt 1998) quote $\alpha = 0.54 \pm 0.04$ and $m_0 = 23.2 \pm 0.10$ for $20 \leq m_R \leq 26$. Chiang & Brown (1999) prefer $\alpha = 0.52 \pm 0.02$ and $m_0 = 23.5$ for $20 \leq m_R \leq 27$; they note that the slope depends on which survey data are used in the fit.

### Table 3

| $M_0/M_{\text{MIN}}$ | $e_0$ | $R_{in}$ (AU) | $R_{out}$ (AU) | $\alpha$ | $m_0$ |
|----------------------|-------|---------------|----------------|----------|-------|
| 0.3 ... ... 10^{-3}  | 42    | 50            | 0.56 $\pm$ 0.01 | 24.03 $\pm$ 0.16 |
| 1.0 ... ... 10^{-4}  | 42    | 50            | 0.58 $\pm$ 0.01 | 23.36 $\pm$ 0.12 |
| 1.0 ... ... 10^{-3}  | 42    | 50            | 0.57 $\pm$ 0.02 | 23.16 $\pm$ 0.15 |
| 1.0 ... ... 10^{-2}  | 42    | 50            | 0.60 $\pm$ 0.01 | 23.53 $\pm$ 0.18 |
| 3.0 ... ... 10^{-3}  | 42    | 50            | 0.58 $\pm$ 0.01 | 22.42 $\pm$ 0.17 |
| 0.3 ... ... 10^{-3}  | 42    | 60            | 0.56 $\pm$ 0.01 | 23.50 $\pm$ 0.13 |
| 1.0 ... ... 10^{-4}  | 42    | 60            | 0.57 $\pm$ 0.01 | 22.85 $\pm$ 0.11 |
| 1.0 ... ... 10^{-3}  | 42    | 60            | 0.56 $\pm$ 0.02 | 22.63 $\pm$ 0.13 |
| 1.0 ... ... 10^{-2}  | 42    | 60            | 0.63 $\pm$ 0.02 | 23.31 $\pm$ 0.19 |
| 3.0 ... ... 10^{-3}  | 42    | 60            | 0.58 $\pm$ 0.01 | 21.92 $\pm$ 0.17 |
constraints on $e_0$. Finally, the proposed Next Generation Space Telescope (NGST) will probe the size distribution of 1 km radius KBOs where models with $e_0 \approx 10^{-3}$ predict a sharp upturn in the observed LF. If such small bodies can survive for the age of the solar system, NGST observations would provide important constraints on the initial mass and dynamics of the outer solar system.

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