Universality in the synchronization of weighted random networks

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(Dated: July 23, 2018)

Over the past few years, the analysis of complex systems from the viewpoint of networks has become an important interdisciplinary issue. It has been shown that physical and dynamical processes, such as cascading failures, epidemic spreading, and network synchronization, are strongly influenced by the structure of the underlying network. Previous works on synchronization have focused mainly on the influence of the topology of the connections by assuming that the coupling strength is uniform. However, synchronization is influenced not only by the topology, but also by the strength of the connections. Most complex networks where synchronization is relevant are indeed weighted. Examples include brain networks, networks of coupled populations in the synchronization of epidemic outbreaks, and technological networks whose functioning relies on the synchronization of interacting units. The distribution of connection weights in real networks is often highly heterogeneous. The study of synchronization in weighted networks is thus of substantial interest.

In this Letter, we address this question in random networks with weighted coupling schemes motivated by real networks. Our main result is the uncovering of a universal formula that describes with good approximation the synchronizability of identical oscillators solely in terms of the mean degree and the heterogeneity of the node’s intensity, irrespective of the degree distribution and other topological properties. The intensity of a node, defined as the total strength of input connections of that node, incorporates both topological and weighted distribution of node’s intensity, where the intensity of a node, denoted as the total strength of input connections, is a natural combination of topology and weights. Our results provide a possibility for the control of synchronization in complex networks by the manipulation of few parameters.

The dynamics of a general weighted network of \( N \) coupled identical oscillators is described by

\[
\dot{x}_i = F(x_i) + \sigma \sum_{j=1}^{N} W_{ij} A_{ij} [H(x_j) - H(x_i)], \quad (1)
\]

\[
\dot{x}_i = F(x_i) - \sum_{j=1}^{N} G_{ij} H(x_j), \quad i = 1, \ldots, N, \quad (2)
\]

where \( F = F(x) \) governs the dynamics of each individual oscillator, \( H = H(x) \) is the output function, and \( \sigma \) is the overall coupling strength. Here \( G = (G_{ij}) \) is the coupling matrix combining both topology [adjacency matrix \( A = (A_{ij}) \)] and weights [weight matrix \( W = (W_{ij}), \ W_{ij} \geq 0 \): \( G_{ij} = \delta_{ij} S_i - W_{ij} A_{ij} \)], where \( S_i = \sum_{j=1}^{N} W_{ij} A_{ij} \) denotes the intensity of node \( i \). The rows of \( G \) have zero sum, and this ensures that the completely synchronized state \( \{x_i = s, \forall i \mid |s| = |F(s)| \} \) is an invariant manifold of Eq. 2. In this work, we focus on the class of weighted networks where \( G \) is diagonalizable and has real eigenvalues. As it will be shown shortly, in this case the synchronizability of the networks can be characterized by the properties of the eigenvalues, without referring to specific forms of \( F \) and \( H \). This applies in particular to the important class of networks where \( G \) can be written as \( G = BC \) for \( B \) a nonsingular diagonal matrix and \( C \) a symmetric, zero row-sum matrix. We assume that matrix \( A \) is binary and symmetric, and that the connection weights and asymmetries are incorporated into \( W \).

The linear stability of the synchronized states can be studied by diagonalizing the variational equations of Eq. 2 into \( N \) blocks of the form

\[
\dot{\xi}_l = [DF(s) - \lambda_l DH(s)]\xi_l, \quad l = 1, \ldots, N, \quad (3)
\]

which are different only by \( \lambda_l \), the \( l \)th eigenvalue of \( G \), ordered as \( 0 = \lambda_1 \leq \lambda_2 \leq \cdots \leq \lambda_N \). The eigenvalues of \( G \) are nonnegative because \( G_{ii} = -\sum_{j \neq i} G_{ij} \geq 0 \) and \( \lambda_1 = 0 \) because \( \sum_{j} G_{ij} = 0 \) for all \( i \). The stability of \( s \) is determined by \( \lambda_1 \) and the master stability function \( \Lambda \), i.e., the largest Lyapunov exponent \( \Lambda \) of the generic variational equation \( \dot{\xi} = [DF(s) - \epsilon DH(s)]\xi \). For many oscillatory dynamical systems, \( \Lambda \) is negative in a
Moreover, close to the synchronized state field of the network, $\bar{\lambda}$ is statistically uncorrelated and that $\Lambda(\sigma \lambda_l) < 0$ for all $l \geq 2$. This is equivalent to the condition
\begin{equation}
R \equiv \lambda_N / \lambda_2 < \epsilon_2 / \epsilon_1,
\end{equation}
where the eigenratio $R$ depends only on the network structure ($G$), and $\epsilon_2 / \epsilon_1$ depends only on the dynamics. From these, it follows that the smaller the eigenratio $R$ the more synchronizable the network and vice versa.

Another measure of synchronizability is the cost $C$ involved in the couplings of the network. When Eq. (4) is satisfied, the synchronized state is linearly stable for $\sigma > \sigma_{\text{min}} \equiv \epsilon_1 / \lambda_2$. The cost $C$ is the total input strength of the connections of all nodes at the synchronization threshold: $C = \sigma_{\text{min}} \sum_{i,j} W_{ij} A_{ij} = \sigma_{\text{min}} \sum_{i=1}^{N} S_i$. The normalized cost\footnote{\[C_0 \equiv C/(N \epsilon_1) = \Omega / \lambda_2,\]}
\begin{equation}
\end{equation}
where $\Omega = \sum_{i=1}^{N} S_i / N$, does not depend on the dynamics ($F$, $H$, and $s$) and can be used as a complementary parameter of synchronizability. Here we characterize the synchronizability of the networks using both the eigenratio $R$ and the cost $C_0$.

Previous work has obtained the eigenvalues of unweighted networks ($W_{ij} = 1$) (see Ref. [15] for a review). Such bounds, however, are not tight and may provide limited information about the actual synchronizability of complex networks. Here we aim at obtaining a more quantitative approximation of the synchronizability for a class of weighted random networks with real spectra, which includes unweighted networks as a special case. Our analysis is based on the combination of a mean field approximation and new graph spectral results.

First, in random networks with $k_{\text{min}} \gg 1$, close to a synchronized state Eq. (1) can be approximated as
\begin{equation}
\dot{x}_i = F(x_i) + \sigma (S_i / k_i) \sum_{j=1}^{N} A_{ij} [H(x_j) - H(x_i)].
\end{equation}
The reason is that each oscillator $j$ receives signals from a large and random sample of other oscillators in the network and $x_j$ is not affected directly by the individual output weights $W_{ij}$. Consequently, we may assume that $W_{ij}$ and $H(x_j)$ are statistically uncorrelated and that $\sum_{j=1}^{N} W_{ij} A_{ij} H(x_j) \approx (1 / k_i) \sum_{j=1}^{N} W_{ij} A_{ij} \sum_{k=1}^{N} A_{ij} H(x_j) = S_i \bar{H}_i$ if $k_i \gg 1$. Here $\bar{H}_i = (1 / k_i) \sum_{j=1}^{N} A_{ij} H(x_j)$ is the local mean field.

Now, if the network is sufficiently random, the local mean field $\bar{H}_i$ can be approximated by the global mean field of the network, $\bar{H}_i \approx H = (1 / N) \sum_{j=1}^{N} H(x_j)$. Moreover, close to the synchronized state $s$, we may assume $\bar{H}_i \approx H(s)$, and the system is approximated as
\begin{equation}
\dot{x}_i = F(x_i) + \sigma S_i [H(s) - H(x_i)], \quad i = 1, \ldots, N,
\end{equation}
indicating that the oscillators are decoupled and forced by a common oscillator $s = F(s)$ with forcing strength proportional to the intensity $S_i$. The variational equations of Eq. (7) have the same form of Eq. (8), except that $\lambda_l$ is replaced by $S_i$. If there exists some $\sigma$ satisfying $\epsilon_1 < \sigma S_i < \epsilon_2$ for all $l$, then all the oscillators are synchronizable by the common driving $H(s)$, corresponding to a complete synchronization of the whole network. These observations suggest that the eigenratio and the cost can be approximated as
\begin{equation}
R \approx S_{\text{max}} / S_{\text{min}}, \quad C_0 \approx \Omega / S_{\text{min}},
\end{equation}
where $S_{\text{min}}, S_{\text{max}}, \text{and } \Omega$ are the minimum, maximum and mean intensities, respectively.

Next we present tight bounds for the above approximation. Eq. (10) means that the coupling matrix $G$ is replaced by the new matrix $G^o = (G^o_{ij})$, with $G^o_{ij} = \delta_{ij} k_i - A_{ij})$. $G^o$ can be written as $G^o = S^o G = SD^{-1}(D - A)$, where $S = (\delta_{ij} S_i)$ and $D = (\delta_{ij} k_i)$ are the diagonal matrices of intensities and degrees, respectively, and $G$ is the normalized Laplacian matrix. Importantly, now the contributions from the topology and weight structure are separated and accounted for $G$ and $S$, respectively. We can show that the largest and smallest nonzero eigenvalues of matrix $G^o$ are bounded by the eigenvalues $\mu_i$ of $G$ as
\begin{equation}
\sigma_{\min} \mu_2 c \leq \lambda_2 \leq \sigma_{\min} c', \quad \sigma_{\max} \leq \lambda_N \leq \sigma_{\max} \mu_N,
\end{equation}
where $c$ and $c'$ can be approximated by 1 for most large complex networks of interest, such as realistic scale-free networks (SFNs). The proof is involved and long, and the details will be presented elsewhere. If the network is sufficiently random, the spectrum of $G$ tends to the semicircle law for large networks with arbitrary expected degrees, provided that $k_{\text{min}} \gg \sqrt{K}$, and $\sigma_{\max} (1 - \mu_2, \mu_N - 1) = [1 + o(1)] \sqrt{K}$ for $k_{\text{min}} \gg \sqrt{K} \ln^3 N$, where $K$ is the mean degree. From these, it follows that
\begin{equation}
\mu_2 \approx 1 - 2 / \sqrt{K}, \quad \mu_N \approx 1 + 2 / \sqrt{K},
\end{equation}
which we find to provide a good approximation under the weaker condition $k_{\text{min}} \gg 1$, regardless of the degree distribution. From Eqs. (10) and (11), we have the following approximations for the bounds of $R$ and $C_0$:
\begin{equation}
\frac{S_{\text{max}}}{S_{\text{min}}} \leq R \leq \frac{S_{\text{max}}}{S_{\text{min}}} [1 + 2 / \sqrt{K}],
\end{equation}
\begin{equation}
\frac{\Omega}{S_{\text{min}}} \leq R \leq \frac{\Omega}{S_{\text{min}}} [1 - 2 / \sqrt{K}].
\end{equation}
For the case of unweighted networks ($S_i = k_i$, $\Omega = K$), the bounds in Eq. (11) are much tighter than those reviewed in Ref. [15].

The bounds in Eqs. (11) and (12) show that the contribution of the network topology is mainly accounted for...
by the mean degree $K$. Therefore, for a given $K$, the synchronizability of random networks with large $k_{\min}$ is expected to be well approximated by the following universal formula:

$$R = A_R \frac{S_{\max}}{S_{\min}}, \quad C_0 = A_C \frac{\Omega}{S_{\min}}, \quad (13)$$

where the pre-factors $A_R$ and $A_C$ are expected to be close to 1. In the case of matrix $G^a$ with uniform intensity ($S_i = 1 \forall i$), they are given by the upper bounds, $A_R = \frac{1+2/\sqrt{K}}{2}$ and $A_C = \frac{1}{2-\sqrt{K}}$, and $A_R \to 1$ and $A_C \to 1$ in the limit $K \to \infty$. Formula (13) is consistent with the approximation in Eq. (9) and indicates that the synchronizability of these networks is primarily determined by the heterogeneity of the intensities.

Our numerical results on various weighted and unweighted networks have confirmed this universal formula. First we consider the following weighted coupling scheme:

$$W_{ij} = S_i / k_i, \quad (14)$$

in which the intensities $S_i$ follow an arbitrary distribution not necessarily correlated with the degrees. In this case, Eq. (9) and Eq. (11) are identical and $G^a = G$. This weighted coupling scheme includes many previously studied systems as special cases. If $S_i = k_i \forall i$, it corresponds to the widely studied case of unweighted networks [2, 4, 5]. In the case of a fully uniform intensity ($S_i = 1 \forall i$), it accommodates a number of previous studies about synchronization of coupled maps [4, 7]. The weighted scheme studied in [2], $W_{ij} = k_i^\theta$, is another special case of Eq. (14) where $S_i = k_i^{1+\theta}$.

We have applied the weighted scheme to various network models: (i) Growing SFNs with aging [8]. Starting with $2m+1$ fully connected nodes, at each time step we connect a new node to $m$ existing nodes according to the probability $P_i \sim k_i^\tau_i$, where $\tau_i$ is the age of the node. The minimum degree is then $k_{\min} = m$ and the mean degree is $K = 2m$. For the aging exponent $\tau \lesssim 0$, this growing exponent generates SFNs with a power-law tail $P(k) \sim k^{-\gamma}$ and scaling exponent in the interval $2 < \gamma \leq 3$ [13], as in most real SFNs. For $\tau = 0$, we recover the usual BA model [13], which has $\gamma = 3$. (ii) Random SFNs [24]. Each node is assigned to have a number $k_i \geq k_{\min}$ of “half-links” according to the distribution $P(k) \sim k^{-\gamma}$. The network is generated by randomly connecting these half-links to form links, prohibiting self- and repeated links. (iii) $K$-regular random networks. Each node is randomly connected to $K$ other nodes.

We present results for two different distributions of intensity $S_i$ which are uncorrelated with the distribution of degree $k_i$: (1) a uniform distribution in $[S_{\min}, S_{\max}]$; and (2) a power-law distribution, $P(S) \sim S^{-\Gamma}$, $S \geq S_{\min}$, where $S_{\min}$ is a positive number. Consistently with the prediction of the universal formula, if $k_{\min} \gg 1$, the eigenratio $R$ collapses into a single curve for a given $K$ when plotted as a function of $S_{\max}/S_{\min}$ [Fig. 1(a)], irrespective of the distributions of $k_i$ and $S_i$. The same happens for the cost $C_0$ as a function of $\Omega/S_{\min}$ [Fig. 1(b)]. The behavior of the fitting parameters $A_R$ and $A_C$ is shown in the insets of Fig. 1. For uniform intensity, they are very close to the upper bounds. They approach very quickly 1 when the intensities become more heterogeneous ($S_{\max}/S_{\min} > 3$). Therefore, Eqs. (13) with $A_R = A_C = 1$ [Fig. 1, solid lines] provide a good approximation of the synchronizability for any large $K$ if the intensities are not very homogeneous.

In more realistic networks, including scientific collaboration networks [12], metabolic networks [21], and airport networks [13, 24], it has been shown that the weight $W_{ij}$ of a connection between nodes $i$ and $j$ is strongly correlated with the product of the corresponding degrees as $\langle W_{ij} \rangle \sim \langle k_i k_j \rangle^\theta$. Here $\theta$ depends on the specific network under study. Motivated by these observations, we analyze the weighted coupling [21].

$$W_{ij} = \langle k_i k_j \rangle^\theta, \quad (15)$$

where the weights are defined for the connections of a given network topology and $\theta$ is a tunable parameter. $\theta$ controls the heterogeneity of the intensity $S_i$ and the correlation between $S_i$ and $k_i$, since $S_i = k_i^{1+\theta}(k_i^\theta)$, where $\langle k_i^\theta \rangle \sim (1/k_i) \sum j k_j^\theta$ is approximately constant for $k_i \gg 1$ when the degree correlations can be neglected. Variations of $\theta$ have significant impact on the synchronizability of SFNs [Fig. 2]. However, as shown in the insets of Fig. 2 for various networks and $\theta$ values, $R$ and $C_0$ collapse again to the universal curves when regarded as functions.
We emphasize that these results also hold for unweighted random networks. As shown in Fig. 3 for unweighted SFNs, the eigenratio $R$ as a function of the system size $N$ collapses to a single universal curve when plotted against $S_{\text{max}}/S_{\text{min}}$.

In addition to heterogeneous degrees and weights, many real networks also display high clustering and nontrivial correlation of degrees. The clustering and degree correlations are negligible in random SFNs and $K$-regular random networks, but we find that they are significant in growing SFNs for $\alpha < 0$. The results in Figs. 1-3 show that the universality holds without significant dependences on these topological properties.

Eqs. (13) also provide a meaningful approximation for networks which are not fully random. For example, consider small-world networks where a regular ring of $N (= 2^{10})$ nodes, each connected to $K (= 20)$ nearest neighbors, is rewired with a probability $p$ for each link. We find that $R$ and $C_0$ collapse to the universal curves even when the networks are dominated by local connections, e.g. for $p = 0.3$, if the intensities are very heterogeneous ($S_{\text{max}}/S_{\text{min}} \geq 10$). For networks with $k_{\text{min}} \sim 1$, the synchronizability is still strongly dependent on $S_{\text{max}}/S_{\text{min}}$ and $K$, although it shows additional dependences on the details of the distributions of $S_i$ and $k_i$ and on other topological properties.

In summary, we have shown that the synchronizability of sufficiently random networks with minimum degree $k_{\text{min}} \gg 1$ is universally dominated by the mean degree $K$ and the heterogeneity of the intensities $S_i$. This universality applies to a general class of large networks where the heterogeneity of $S_i$ is due to either the distribution of degrees, as in unweighted SFNs, or the distribution of connection weights, as in weighted $K$-regular networks, or a combination of both, as expected in most realistic networks, such as in the airport network which underlies the synchronization of epidemic outbreaks. In particular, formula (13) explains why synchronizability is improved when the heterogeneity of $S_i$ is reduced, which can be useful for network design and control of synchronization.

C. S. Z. and J. K. were supported by the VW Foundation and SFB 555. A. E. M. was supported by DOE under Contact No. W-7405-ENG-36.

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