On the multifractal local behavior of parabolic stochastic PDEs

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Abstract

Consider the stochastic heat equation \( \dot{u} = \frac{1}{2} u'' + \sigma(u) \xi \) on \((0, \infty) \times \mathbb{R}\) subject to \( u(0) \equiv 1 \), where \( \sigma : \mathbb{R} \to \mathbb{R} \) is a Lipschitz (local) function that does not vanish at 1, and \( \xi \) denotes space-time white noise. It is well known that \( u \) has continuous sample functions \([22]\); as a result, \( \lim_{t \downarrow 0} u(t, x) = 1 \) almost surely for every \( x \in \mathbb{R} \).

The corresponding fluctuations are also known \([14, 16, 20]\): For every fixed \( x \in \mathbb{R} \), \( t \mapsto u(t, x) \) looks locally like a fixed multiple of fractional Brownian motion (fBm) with index \( 1/4 \). In particular, an application of Fubini’s theorem implies that, on an \( x \)-set of full Lebesgue measure, the short-time behavior of the peaks of the random function \( t \mapsto u(t, x) \) are governed by the law of the iterated logarithm for fBm, up to possibly a suitable normalization constant. By contrast, the main result of this paper claims that, on an \( x \)-set of full Hausdorff dimension, the short-time peaks of \( t \mapsto u(t, x) \) follow a non-iterated logarithm law, and that those peaks contain a rich multifractal structure a.s.

Large-time variations of these results were predicted in the physics literature a number of years ago and proved very recently in \([10, 11]\). To the best of our knowledge, the short-time results of the present paper are observed here for the first time.

Keywords: the stochastic heat equation; multifractals; Hausdorff dimension; packing dimension.

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1 Introduction

Let \( \sigma : \mathbb{R} \to \mathbb{R} \) be a non-random, Lipschitz continuous function, \( \xi := \) space-time white noise, and consider the unique, continuous solution \( u = u(t, x) \) to the semi-linear stochastic heat equation,

\[
\begin{cases}
\dot{u} = \frac{1}{2} u'' + \sigma(u) \xi & \text{on } (0, \infty) \times \mathbb{R}, \\
\text{subject to } u(0) \equiv 1.
\end{cases}
\]

(1.1)

It is not hard to see that if \( \sigma(1) = 0 \), then \( u \equiv 1 \) a.s. Therefore, we assume that

\( \sigma(1) \neq 0, \)

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to avoid trivialities. In this case, it has been shown recently by Khoshnevisan, Swanson, Xiao, and Zhang [14] that, for every fixed \( x \in \mathbb{R} \) and \( t_0 > 0 \), there exists a fractional Brownian motion \( B = \{ B_t \}_{t \geq 0} \) of index \( 1/4 \) such that

\[
    u(t + t_0, x) - u(t_0, x) \approx (2/\pi)^{1/4} \sigma(u(t_0, x)) B_t \quad \text{a.s. when } t \approx 0.
\]

For closely-related works, see also Lei and Nualart [16] and Pospíšil and Tribe [20].

Throughout, \( \dim_H \) and \( \dim_p \), respectively denote the Hausdorff and the packing dimension; see Matilla [18, Chs 4 and 5]. We will also let \( \dim_M \) denote the upper Minkowski—or box—dimension; see the book by Matilla (ibid.).

**Theorem 1.1.** If \( 0 < c \leq \sigma(1)/\pi^{1/4} \), then with probability one,

\[
    \dim_p(\mathcal{U}(c)) = 1 \quad \text{and} \quad \dim_M(\mathcal{U}(c)) = 1 - \frac{c^2 \sqrt{\pi}}{\| \sigma(1) \|^2}.
\]

If \( c > \sigma(1)/\pi^{1/4} \), then \( \mathcal{U}(c) = \emptyset \) a.s.

It has recently been shown by Kunwoo Kim, Yimin Xiao, and the second author [10, 11] that the largest global oscillations of the random map \((t, x) \mapsto u(t, x)\) form an asymptotic multifractal; this property had been predicted earlier in a voluminous physics literature on the present “intermittent” stochastic systems [as well as for more complex systems]. It is not hard to deduce from Theorem 1.1 that the high local oscillations of \((t, x) \mapsto u(t, x)\) are also multifractal in a precise sense that we now explain next. Indeed, let us observe that, by Theorem 1.1,

\[
    \mathcal{U} = \left\{ x \in [0, 1] : 0 < \limsup_{\varepsilon \downarrow 0} \frac{u(\varepsilon, x) - 1}{\varepsilon^{1/4} \sqrt{\log(1/\varepsilon)}} < \infty \right\} \quad \text{a.s.}
\]

Moreover, the right-hand side of (1.4) is equal to \( \cup_{c \in I} \mathcal{U}(c) \)—where \( I := (0, \sigma(1)/\pi^{1/4}] \)—and \( c \mapsto \dim_M(\mathcal{U}(c)) \) is strictly decreasing on \( I \). Thus, Theorem 1.1 implies that the
right-hand side of (1.4) is a “multifractal.”

Finally, let us recall that by a “solution” \( u = u(t, x) \) to equation (1.1) we mean that \( u \) is a “mild” solution to (1.1). That is, \( u \) satisfies the following for each \( t > 0 \) and \( x \in \mathbb{R} \):

\[
    u(t, x) = 1 + \int_{(0,t) \times \mathbb{R}} p_{t-s}(x-y) \sigma(u(s, y)) \xi(dy) \quad \text{a.s.,} \tag{1.5}
\]

where \( p_r(w) := (2\pi r)^{-1/2} e^{-|w|^2/(2r)} \) \([r > 0, w \in \mathbb{R}]\) denotes the heat kernel on \( \mathbb{R} \), and the stochastic integral is defined by Walsh [22].

2 The constant-coefficient case

Consider first the constant-coefficient stochastic heat equation,

\[
    \begin{cases}
    \dot{Z} = \frac{1}{2} Z'' + \xi & \text{on } (0, \infty) \times \mathbb{R}, \\
    \text{subject to } Z(0) = 0.
    \end{cases}
\]  

It is easy to see that \( Z(t, x) = u(t, x) - 1 \), where \( u \) solves (1.1) in the special case that \( \sigma \equiv 1 \).

According to the theory of Walsh [22, Ch. 3], the solution to (2.1) can be written, in mild form, as the following Wiener integral process:

\[
    Z(t, x) = \int_{(0,t) \times \mathbb{R}} p_{t-s}(y-x) \xi(dy).
\]  

Evidently, \( Z \) is a centered, Gaussian random field. The following simple computation contains all of the requisite information about the process \( Z \).

**Lemma 2.1.** For all \( \varepsilon > 0 \) and \( x, x' \in \mathbb{R} \),

\[
    \operatorname{Cov}[Z(\varepsilon, x), Z(\varepsilon, x')] = \frac{\varepsilon}{2} g_2 \left( \frac{x-x'}{\sqrt{\varepsilon}} \right),
\]

where \( g \) denotes the incomplete Green’s function of the heat kernel; that is, \( g_t(a) := \int_0^t p_r(a) \, dr \) for all \( t > 0 \) and \( a \in \mathbb{R} \).

**Proof.** Let \( \delta := x - x' \). By the Wiener isometry and (2.2),

\[
    \operatorname{Cov}[Z(\varepsilon, x), Z(\varepsilon, x')] = \int_0^\varepsilon ds \int_{-\infty}^\infty dy p_{t-s}(y-x)p_{t-s}(y-x') = \int_0^\varepsilon p_{2s}(\delta) \, ds.
\]

We have used the semigroup property of the heat kernel in the second identity. We can change variables \(|r := 2s/\varepsilon|\) and observe that \( p_{\varepsilon \delta}(\delta) = \varepsilon^{-1/2} p_{\delta}(\delta/\sqrt{\varepsilon}) \) to finish. \( \square \)

By l’Hôpital’s rule, \( g_2(a) \sim 4\pi^{-1/2} a^{-2} \exp(-a^2/4) \) as \( a \to \infty \). Therefore, Lemma 2.1, and the strict positivity and the continuity of \( g_2 \) together imply the following: The Gaussian random field \( Z(\varepsilon) \) is stationary for every fixed \( \varepsilon > 0 \). Furthermore, for every real number \( \alpha \),

\[
    \operatorname{Cov}\left[Z(\varepsilon, 0), Z\left(\varepsilon, \varepsilon^{(1/2) + \alpha}\right)\right] \sim \begin{cases} 
    \varepsilon^{1/2} & \text{if } \alpha \geq 0, \\
    \varepsilon^{(1/2) + 2|\alpha|} \exp\left(-\frac{1}{4\varepsilon^{2|\alpha|}}\right) & \text{if } \alpha < 0,
    \end{cases}
\]

uniformly for all \( \varepsilon \in (0, 1) \). This implies roughly that if \(|x-y| \gg \sqrt{\varepsilon}\) then \( Z(\varepsilon, x) \) and \( Z(\varepsilon, y) \) are very close to being uncorrelated, whereas \( Z(\varepsilon) \) locally behaves as a random
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constant on a spatial scale of $O(\sqrt{\varepsilon})$. In other words, the “correlation length” of the random field $Z(\varepsilon)$ is $\sqrt{\varepsilon}$ when $\varepsilon \approx 0$.

The rest of this section is devoted to developing some hitting estimates for $Z$. The latter is a far simpler object than the solution $u$ to (1.1) when $\sigma$ is a non-linear function. Therefore, one can hope for better information about $Z$ than $u$.

2.1 Upper bounds

Because $Z$ is a Gaussian random field, we can appeal to concentration of measure ideas in order to bound the local supremum of $Z$. The following entropy estimate contains the key step in that direction.

**Lemma 2.2.** Choose and fix a real number $R > 0$, and define

$$B_R(\varepsilon) := (0, \varepsilon) \times [0, R\sqrt{\varepsilon}] \quad \text{for all } \varepsilon > 0. \quad (2.3)$$

Then, $E[\sup_{B_R(\varepsilon)} Z] \lesssim \varepsilon^{1/4}$, uniformly for all $\varepsilon \in (0, 1)$.

**Proof.** It is well known that $\sqrt{E[(Z(t, x) - Z(s, y))^2]} \lesssim d((t, x); (s, y))$, uniformly for all $s, t \geq 0$ and $x, y \in \mathbb{R}$, where $d$ denotes the spatio-temporal distance function,

$$d((t, x); (s, y)) := |t - s|^{1/4} + |x - y|^{1/2}. \quad (2.4)$$

This is a well-known part of the folklore of the subject, and appears upon close inspection within the proof of Corollary 3.4 of Walsh [22, p. 318], for example. The details are worked out pedagogically in [2, (135), p. 31].

Next we note that by Dudley’s theorem [3],

$$E \sup_{B_R(\varepsilon)} Z \lesssim \int_0^{D(\varepsilon)} \sqrt{\log N_\varepsilon(r)} \, dr, \quad (2.5)$$

uniformly for all $\varepsilon \in (0, 1)$, where:

(a) $D(\varepsilon)$ denotes the $d$-diameter of $B_R(\varepsilon)$;

(b) $N_\varepsilon$ is the metric entropy of $B_R(\varepsilon)$; that is, $N_\varepsilon(r)$ denotes the minimum number of $d$-balls of radius $r > 0$ that are needed to cover $B_R(\varepsilon)$; and

(c) The implied constant in the inequality (2.5) does not depend on $\varepsilon$ (consult, for example, Marcus and Rosen [17, Theorem 6.2, p. 245]).

We may consider covering $B_R(\varepsilon)$ by $d$-balls of the form $\{(s, y) : |s - t|^{1/4} + |y - x|^{1/2} \leq r\}$ for every fixed $t, r > 0$ and $x \in \mathbb{R}$. Clearly, the number of such $d$-balls of radius $r > 0$ that are needed to cover $B_R(\varepsilon)$ is of sharp order

$$(1 + \frac{\varepsilon}{r^4}) \left(1 + \frac{\sqrt{\varepsilon}}{r^2}\right) \approx \left(1 + \frac{\varepsilon}{r^4}\right)^{3/2} := M_\varepsilon(r),$$

uniformly for all $\varepsilon > 0$ and $r \in (0, D(\varepsilon))$. It is not hard to see that $D(\varepsilon) \gtrsim \varepsilon^{1/4}$, uniformly for all $\varepsilon \in (0, 1)$. Because $N_\varepsilon \lesssim M_\varepsilon$ uniformly on $(0, D(\varepsilon))$, it follows from (2.5) that there exists a real number $C > 0$ such that

$$E \sup_{B_R(\varepsilon)} Z \lesssim \int_0^{C\varepsilon^{1/4}} \sqrt{1 + \log \left(1 + \frac{\varepsilon}{r^4}\right)} \, dr \propto \varepsilon^{1/4},$$

uniformly for all $\varepsilon \in (0, 1)$. This completes the proof. \qed

Next we use Lemma 2.2 and concentration of measure in order to deduce an asymptotically-sharp maximal inequality for $Z$. 

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Lemma 2.3. Choose and fix a real number $R > 0$, and recall (2.3). Then,

$$\lim_{\lambda \to \infty} \frac{1}{\lambda^2} \log P \left\{ \sup_{B_n(\varepsilon)} Z \geq \left( \frac{4\varepsilon}{\pi} \right)^{1/4} \lambda \right\} = -1,$$

uniformly for all $\varepsilon \in (0, 1)$.

Proof. In accord with Lemma 2.1, the Gaussian random field $Z(\varepsilon)$ is stationary and centered, with $\text{Var}[Z(\varepsilon, 0)] = \sqrt{\varepsilon/\pi}$. Therefore, the inequality of Borell [1] and Sudakov and T’ivrelos [21] implies that $\sup_{B_n(\varepsilon)} Z - E\sup_{B_n(\varepsilon)} Z$ has sub-Gaussian tails; see Ledoux [15, Ch. 7]. After a few lines of computations, this fact and Lemma 2.2 together imply that

$$\sup_{\varepsilon \in (0, 1)} \log P \left\{ \sup_{B_n(\varepsilon)} Z \geq \left( \frac{4\varepsilon}{\pi} \right)^{1/4} \lambda \right\} \leq -\lambda^2 + o(\lambda),$$

as $\lambda \to \infty$. At the same time,

$$P \left\{ \sup_{B_n(\varepsilon)} Z \geq \left( \frac{4\varepsilon}{\pi} \right)^{1/4} \lambda \right\} \geq P \left\{ Z(\varepsilon, 0) \geq \left( \frac{4\varepsilon}{\pi} \right)^{1/4} \lambda \right\} = P\{X \geq 2^{1/2}\lambda\},$$

where $X$ has a standard normal distribution; see Lemma 2.1. Since $\log P\{X \geq 2^{1/2}\lambda\} \geq -\lambda^2 + o(\lambda)$, as $\lambda \to \infty$, the result follows. \hfill \Box

Define, for all $c > 0$, the random set

$$\mathcal{G}(c) := \left\{ y \in [0, 1] : \limsup_{\varepsilon \downarrow 0} \frac{Z(\varepsilon, y)}{\varepsilon^{1/4} \sqrt{\log(1/\varepsilon)}} \geq c \right\}. \quad \text{(2.6)}$$

Lemma 2.4. Let $E$ be a measurable subset of $[0, 1]$. Then, $P\{\mathcal{G}(c) \cap E = \emptyset\} = 1$ for every $c > (1/\pi)^{1/4} \sqrt{\dim_{\text{H}}(E)}$.

Proof. Define, for every $c > 0$,

$$\mathcal{L}(c) := \left\{ t \in [0, 1] : \sup_{y \in E} \frac{Z(t, y)}{t^{1/4} \sqrt{\log(1/t)}} \geq c \right\}.$$

Every $\mathcal{L}(c)$ is a random subset of the time interval $[0, 1]$. For all integers $n \geq 1$ define

$$t_n := e^{-\sqrt{n}}, \quad \mathcal{T}(n) := [t_{n+1}, t_n], \quad \text{and} \quad \mathcal{A}(j; n) := [j\sqrt{t_n}, (j+1)\sqrt{t_n}],$$

for all integers $j \geq 0$. Then, by monotonicity and a simple union bound,

$$P\{\mathcal{L}(c) \cap \mathcal{T}(n) \neq \emptyset\} \leq \sum_{0 \leq j < 1/t_n; \mathcal{A}(j; n) \cap E \neq \emptyset} P \left\{ \sup_{t \in \mathcal{T}(n)} \sup_{y \in \mathcal{A}(j; n)} Z(t, y) > c(n t_{n+1})^{1/4} \right\}.$$

Recall the space-time sets $B_n(\varepsilon)$ from (2.3). Since $y \mapsto Z(t, y)$ is stationary, the natural logarithm of the $j$th term of the preceding sum is equal to

$$\log P \left\{ \sup_{t \in \mathcal{T}(n)} \sup_{y \in \mathcal{A}(j; n)} Z(t, y) > c(n t_{n+1})^{1/4} \right\} \leq \log P \left\{ \sup_{B_1(\varepsilon)} Z > c(n t_{n+1})^{1/4} \right\} \sim -\frac{c^2 \sqrt{\pi n}}{2},$$
as \( n \to \infty \), thanks to Lemma 2.3 and the fact that \( t_{n+1}/t_n = 1 - (\frac{1}{2} + o(1))n^{-1/2} \) as \( n \to \infty \). Therefore, by the definition of the Minkowski dimension,

\[
\limsup_{n \to \infty} \frac{\log P \{ \mathcal{L}(c) \cap \mathcal{T}(n) \neq \emptyset \}}{\sqrt{n}} \leq -\frac{c^2\sqrt{\pi}}{2} + \frac{\dim_p(E)}{2}.
\]

In particular, if the right-hand side is \(< 0\), then \( n \to P \{ \mathcal{L}(c) \cap \mathcal{T}(n) \neq \emptyset \} \) sums to a finite number. Therefore, the Borel–Cantelli lemma implies that

\[
\limsup_{n \to \infty} (\mathcal{L}(c) \cap \mathcal{T}(n)) = \emptyset \quad \text{a.s. for all } \forall c > (1/\pi)^{1/4} \sqrt{\dim_p(E)}.
\]

It is easy to see that \( \mathcal{S}(c) \cap E = \emptyset \) a.s. on the event that \( \limsup_{n \to \infty} (\mathcal{L}(c) \cap \mathcal{T}(n)) = \emptyset \), whence follows the lemma. \( \square \)

It is easy to apply Lemma 2.4 in order to improve itself slightly. Note that the difference between the following and Lemma 2.4 is one about the Minkowski versus the packing dimension of \( E \).

**Lemma 2.5.** Let \( E \) be a measurable subset of \([0, 1]\). Then, \( P \{ \mathcal{S}(c) \cap E = \emptyset \} = 1 \) for every \( c > (1/\pi)^{1/4} \sqrt{\dim_p(E)} \).

**Proof.** Choose and fix \( \eta > 0 \) small enough to ensure that \( c > (1/\pi)^{1/4} \sqrt{\dim_p(E)} + \eta \). By the definition of packing dimension [18, p. 81] there exists a closed cover \( F_1, \ldots, F_m \) of \( E \) such that \( \dim_p(E) \geq \max_{1 \leq j \leq m} \dim_p(F_j) - \eta \). Thus, \( \mathcal{S}(c) \cap F_j = \emptyset \) a.s. for all \( 1 \leq j \leq m \) by Lemma 2.4. This completes the proof. \( \square \)

### 2.2 Lower bounds

As was mentioned earlier, the “correlation length” of the spatial process \( Z(\varepsilon) \) is of sharp order \( \sqrt{\varepsilon} \) when \( \varepsilon \) is small. It is possible to state and prove a much stronger, quantitative version of this assertion. In order to describe that, let us define for all \( x \in \mathbb{R} \) and \( \varepsilon, \delta \in (0, 1) \),

\[
\Gamma(x; \varepsilon, \delta) := \left[ x - \sqrt{2\varepsilon \log(1/\delta)}, x + \sqrt{2\varepsilon \log(1/\delta)} \right],
\]

and

\[
Z_\delta(\varepsilon, x) := \int_{\Gamma(x; \varepsilon, \delta)} p_{\varepsilon, s}(y - x) \xi(ds, dy).
\]

**Lemma 2.6.** \( E(\|Z(\varepsilon, x) - Z_\delta(\varepsilon, x)\|^2) < \delta^2 \sqrt{\varepsilon} \) for all \( \varepsilon, \delta \in (0, 1) \) and \( x \in \mathbb{R} \).

**Remark 2.7.** Recall that if \( \phi, \psi \in L^2(\mathbb{R} \times \mathbb{R}) \) are orthogonal in \( L^2(\mathbb{R} \times \mathbb{R}) \) then the Wiener integrals \( \int \phi \, d\xi \) and \( \int \psi \, d\xi \) are independent. This observation has the following by-product: If \( x_1, \ldots, x_k \in \mathbb{R} \) satisfy

\[
\min_{i \neq j} |x_i - x_j| > \sqrt{8 \varepsilon \log(1/\delta)},
\]

then \( \{Z_\delta(\varepsilon, x_i)\}_{i=1}^k \) are independent and identically distributed. It follows that the gap condition (2.9) ensures that \( \{Z(\varepsilon, x_i)\}_{i=1}^k \) is uniformly to within \( O(\delta^{1/4}) \) of an i.i.d. sequence.

**Proof of Lemma 2.6.** By stationarity, we may—and will—consider only the case that \( x = 0 \). According to the Wiener isometry, \( E(\|Z(\varepsilon, 0) - Z_\delta(\varepsilon, 0)\|^2) \) is equal to

\[
\int_0^\varepsilon ds \int_{y \in \mathbb{R}: |y| > \sqrt{2|s| \log \delta}} dy \left| p_s(y) \right|^2 \leq \int_0^\varepsilon P \left\{ |X| > \sqrt{\frac{4|s| \log \delta}{\pi s}} \right\} \frac{ds}{\sqrt{2\pi s}}.
\]

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where $X$ has the standard normal distribution. Elementary manipulations and the well-known tail bound $P\{|X| > \lambda\} \leq \exp(-\lambda^2/2)$ together yield the lemma.

\begin{lemma} \label{lem:tail-bound}
Let $F_1, F_2, \ldots$ be finite subsets of $\mathbb{R}$ that satisfy $\log |F_n| \leq \kappa(1+o(1))n$ as $n \to \infty$, for some $0 \leq \kappa < \infty$. Then, uniformly for all $0 < \delta < 1$ and $0 < \nu < \exp(-2\kappa)$,
\[
\limsup_{n \to \infty} \max_{x \in F_n} \frac{|Z(\nu^n, x) - Z_\delta(\nu^n, x)|}{\nu^{n/4} \sqrt{\log(1/\nu^n)}} \leq \delta \quad \text{a.s.}
\]
\end{lemma}

\begin{proof}
Lemma 2.6 ensures that
\[
\text{Var}[Z(\nu^n, 0) - Z_\delta(\nu^n, 0)] \leq \nu^{n/2} \delta^2 \quad \text{for all } n \geq 1.
\]
Thanks to this and stationarity, a standard Gaussian tail bound yields
\[
\frac{1}{n} \log P \left\{ \max_{x \in F_n} |Z(\nu^n, x) - Z_\delta(\nu^n, x)| > \delta \nu^{n/4} \sqrt{\log(1/\nu^n)} \right\} \leq \frac{1}{2} \log(1/\nu) + \kappa + o(1),
\]
as $n \to \infty$. The right-hand side of the preceding display is strictly negative for all $n$ sufficiently large. Therefore, the Borel–Cantelli implies the result.
\end{proof}

Finally, we derive a matching “converse” to Lemma 2.5. The proof of the following result borrows heavily ideas from the theory of limsup random fractals [12], see also [13].

\begin{lemma} \label{lem:converse}
Recall (2.6), and let $E$ denote an arbitrary measurable subset of $[0, 1]$. Then, $P\{\mathcal{G}(c) \cap E \neq \emptyset\} = 1$ for every strictly positive constant $c < \pi^{-1/4} \sqrt{\dim_o(E)}$.
\end{lemma}

\begin{proof}
Choose and fix a number $\gamma \in (0, \dim_o(E))$. A theorem of Joyce and Preiss [8] implies that there exists a compact set $F \subseteq E$ such that $\dim_o(F) > \gamma$. Define $E_* := \cap_{n=1}^\infty \cup (F \cap I)$, where the union is over all dyadic subintervals of $[0, 1]$ that have length $2^{-n}$ and satisfy $\dim_o(F \cap I) \geq \gamma$. Clearly, $E_* \subseteq E$ is compact. And the definition of packing dimension implies that
\[
\dim_p(E_* \cap I) \geq \gamma \quad \text{for all open intervals } I \text{ that intersect } E_*.
\]

Let $\delta \in (0, 1/(2\pi^{1/2}))$ and $\nu \in (0, 1)$ be fixed. Also, fix an arbitrary open interval $I$ that intersects $E_*$, and let $K$ denote the Kolmogorov capacity—or packing numbers—of the compact set $E_* \cap I$. That is, for every $\varepsilon > 0$, $K(\varepsilon)$ denote the maximum integer $k$ such that we can find $a_1, \ldots, a_k \in E_* \cap I$ with the property that $\min_{1 \leq i \neq j \leq k} |a_i - a_j| > \varepsilon$. Recall [18, p.~78] that $\dim_p(E_* \cap I) = \limsup_{\varepsilon \downarrow 0} \log K(\varepsilon)/\log(1/\varepsilon)$. Because of (2.10), the preceding fact has the following consequence: There exists an infinite collection $\mathcal{N}$ of positive integers, and a sequence $x_{1, n}, \ldots, x_{N(n), n} \in E_* \cap I$—for every $n \geq 1$—with the following properties:

(a) $\min_{1 \leq i \neq j \leq N(n)} |x_{i, n} - x_{j, n}| > \sqrt{8\nu^n \log(1/\delta)}$ for all $n \geq 1$;
(b) There exists $M > 0$ such that $N(n) \leq M [\nu^n \log(1/\delta)]^{-\gamma/2}$ for all $n \geq 1$; and
(c) $N(n) \geq 8\nu^n \log(1/\delta) \gamma/2$ for all $n \in \mathcal{N}$.

By Remark 2.7, and because of the stationarity of $x \mapsto Z_\delta(t, x)$, the random variables
\[
\{Z_\delta(\nu^n, x_{i, n})\}_{i=1}^{N(n)}
\]
are i.i.d.

For every $c > 0$ and $n \in \mathcal{N}$, let
\[
S_n(c) := \sum_{i=1}^{N(n)} I \left\{ Z_\delta(\nu^n, x_{i, n}) > c\nu^{n/4} \sqrt{\log(1/\nu^n)} \right\}.
\]
We have shown that if $c$ where $I$ holds automatically because $\gamma < \dim(E) \leq 1$. Because $\delta$ can be otherwise as small as we wish, and since $\gamma < \dim(E)$ can be as large as we want, this proves the lemma.

Finally, let us describe the critical case that is left open in Lemma 2.9.

**Lemma 2.10.** The conclusion of Lemma 2.9 continues to hold when $c = \pi^{-1/4} \sqrt{\dim(E)}$.

**Proof.** Choose and fix a countable sequence of real numbers $0 < c_1 < c_2 < \cdots$ that converge upward to $c = \pi^{-1/4} \sqrt{\dim(E)}$. The proof of Lemma 2.9 showed—for the same compact set $E_2 \subset E$ as before—that $P[\mathcal{G}(c_1) \cap E_2 \neq \emptyset] = 1$ for all $n \geq 1$ and all open intervals $I$ that intersect $E_2$. An application of the Baire category theorem reveals that $\mathcal{G}(c) = \cap_{n=1}^\infty \mathcal{G}(c_n)$ a.s. intersects $E_2$ and hence $E$.

**3 Proof of Theorem 1.1**

The proof of Theorem 1.1 hinges on a localization result which relates the small-time behavior of $u$ to that of $Z$. First, let us recall a moment estimate of Khoshnevisan, Swanson, Xiao, and Zhang. When (1.1) is replaced by the stochastic heat equation on a torus, an almost-sure version of the following was derived in the seminal work of Hairer [4], and played an important role in the ensuing deep theory of regularity structures of Hairer [5]. An almost-sure version of the following [for (1.1), as is, on the real line] can likely be deduced also from the work of Hairer and Labbé [6].
We may apply Proposition 3.1 [with an arbitrary $k$ where $\Phi$ is in $\{u, Z\}$] uniformly for all $\epsilon > 0$, and the symbol $\Phi$ is in $\{u, Z\}, K(u) = 1$, and $K(Z) = 1/\sigma(1)$. Because $\sigma(1) \neq 0$ throughout, $K(\Phi)$ is well defined and finite. Therefore, by the work of Walsh [22, Ch. 3],

$$
\sup_{t \in [0, \epsilon]} \sup_{x \in [0, 1]} |\Phi(t, x) - \Phi(s, y)| \leq e^{\theta + o(1)} \quad \text{as } \epsilon \downarrow 0,
$$

for both possible choices of $\Phi \in \{u, Z\}$. As a result, it can be deduced from Chebyshev's

**Proposition 3.1** ([14]). For every $k \in [2, \infty)$ and $\theta \in (0, 2/5)$,

$$
\sup_{t \in (0, \epsilon)} \sup_{x \in \mathbb{R}} E \left( |u(t, x) - 1 - \sigma(1)Z(t, x)|^k \right) = o \left( \epsilon^{k\theta} \right) \quad \text{as } \epsilon \downarrow 0.
$$

We now apply Proposition 3.1—with $k = 2$—in order to establish the following.

**Proposition 3.2.** For every $\theta \in (0, 2/5)$ and $N, M \geq 0$,

$$
\sup_{t \in (0, \epsilon)} \sup_{x \in [-M, N]} |u(t, x) - 1 - \sigma(1)Z(t, x)| = o \left( \epsilon^{\theta} \right) \quad \text{a.s.}
$$

**Remark 3.3.** If, instead of an SPDE on $\mathbb{R} \times \mathbb{R}$, we studied an analogous SPDE on $\mathbb{R} \times [0, 1]$, then the work of Hairer and Pardoux [7] improves further the error rate of the analogue of Proposition 3.2 to $O(\sqrt{\epsilon} \log \epsilon)$. We are not sure if the latter rate is the right one in the present setting, but will not need these improvements and so will prove only what we need.

**Proof of Proposition 3.2.** Without loss of great generality, we study only the case that $M = 0$ and $N = 1$; the general case is proved similarly.

Let $\Delta(t, x) := u(t, x) - 1 - \sigma(1)Z(t, x)$ for all $t \geq 0$ and $x \in \mathbb{R}$, and for every $\epsilon \in (0, 1)$ define

$$
\Xi(\epsilon) := \left\{ je^{2\theta} : 0 \leq j < e^{-2\theta}, j \in \mathbb{Z} \right\} \quad \text{and} \quad \Xi(\epsilon) := \left\{ je^{4\theta} : 0 \leq j < e^{1-4\theta}, j \in \mathbb{Z} \right\}.
$$

We may apply Proposition 3.1 [with an arbitrary $k \geq 2$] in order to see that, for all real numbers $b \in (0, \theta)$,

$$
P \left( \max_{t \in \Xi(\epsilon)} \max_{x \in \Xi(\epsilon)} |\Delta(t, x)| > \epsilon^b \right) \leq e^{-kb} \sum_{t \in \Xi(\epsilon)} \sum_{x \in \Xi(\epsilon)} \|\Delta(t, x)\|^k \leq e^{(k-1)\theta - kb + 1},
$$

uniformly for all $\epsilon \in (0, 1)$. Therefore,

$$
P \left( \sup_{t \in (0, \epsilon)} \sup_{x \in [0, 1]} |\Delta(t, x)| > 3\epsilon^b \right) \leq e^{(k-1)\theta - kb + 1} + \mathcal{P}(u; \epsilon) + \mathcal{P}(Z; \epsilon),
$$

where

$$
\mathcal{P}(\Phi; \epsilon) := P \left( \sup_{t \in (0, \epsilon)} \sup_{x \in [0, 1]} |\Phi(t, x) - \Phi(s, y)| > K(\Phi)\epsilon^b \right)
$$

and the symbol $\Phi$ is in $\{u, Z\}, K(u) = 1$, and $K(Z) = 1/\sigma(1)$. Because $\sigma(1) \neq 0$ throughout, $K(\Phi)$ is well defined and finite. Therefore, by the work of Walsh [22, Ch. 3],
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inequality that \( P(\Phi; \varepsilon) \leq \varepsilon^{k(\theta-b)+o(1)} = O(\varepsilon^{k(\theta-b)}) \). Thus, we find that

\[
\begin{align*}
P \left( \sup_{t \in (0, \varepsilon)} \sup_{x \in [0,1]} |\Delta(t,x)| > 3\varepsilon^b \right) \leq \varepsilon^{(k-6)\theta-kb} \quad \text{as} \ \varepsilon \downarrow 0.
\end{align*}
\]

Replace \( \varepsilon \) by \( 2^{-n} \) to see from a monotonicity argument and the Borel-Cantelli lemma that, almost surely, \( \sup_{t \in (0,\varepsilon)} \sup_{x \in [0,1]} |\Delta(t,x)| = O(\varepsilon^b) \) as \( \varepsilon \downarrow 0 \). The result follows from this because \( k \geq 2, \theta \in (0, 2/5), b \in (0, (k-6)\theta/k) \) are all otherwise arbitrary.

We conclude the paper with the following.

Proof of Theorem 1.1. We will assume throughout that \( E \) is a subset of \( [0,1] \). A routine change of scale [in the arguments] will yield the general case.

Proposition 3.2 implies that, for every real number \( c > 0 \), \( \mathcal{W}(c) = \mathcal{W}(c/\sigma(1)) \) a.s., where the random sets \( \mathcal{G}(\bullet) \) were defined in (2.6). It follows from this and Lemmas 2.5, 2.9, and 2.10 that for all measurable sets \( E \subset [0,1] \):

A1. If \( c \leq \pi^{-1/4}\sigma(1)\sqrt{\dim_\nu(E)} \), then \( P(\mathcal{W}(c) \cap E \neq \emptyset) = 1 \); and

A2. If \( c > (1/\pi)^{1/4}\sigma(1)\sqrt{\dim_\nu(E)} \), then \( P(\mathcal{W}(c) \cap E \neq \emptyset) = 0 \); and

According to the theory of limsup random fractals [12], for every \( \rho \in (0,1) \) there exists a random set \( \Sigma_\rho \subset [0,1] \), that is independent of the process \( u \), and satisfies the following for all measurable sets \( F \subset [0,1] \):

B1. \( P(\Sigma_\rho \cap F \neq \emptyset) = 1 \) if \( \rho < \dim_\nu(F) \). In this case, \( \dim_\nu(\Sigma_\rho \cap F) = \dim_\nu(F) \); and

B2. \( P(\Sigma_\rho \cap F \neq \emptyset) = 0 \) if \( \rho > \dim_\nu(F) \).

Both B1 and B2 follow from Theorems 3.1 and 3.2 of [12].

By the independence of \( \Sigma_\rho \) and \( \mathcal{W}(c) \), we may combine A1 and B1-B2 with \( F = \mathcal{W}(c) \) and \( E = \Sigma_\rho \), in order to see that \( \dim_\nu(\Sigma_\rho) = 1 \) and:

\[
c \leq \sigma(1)/\pi^{1/4} \Rightarrow P(\mathcal{W}(c) \cap \Sigma_\rho \neq \emptyset) = 1 \Rightarrow \rho \leq \dim_\nu(\mathcal{W}(c)).
\]

Since \( \rho \in (0,1) \) was arbitrary, it follows that \( \dim_\nu(\mathcal{W}(c)) = 1 \) a.s. for all \( c \leq \sigma(1)/\pi^{1/4} \).

On the other hand, A2 shows that if \( c > \sigma(1)/\pi^{1/4} \), then \( \mathcal{W}(c) = \emptyset \) a.s.

Finally, the computation of the Hausdorff dimension of \( \mathcal{W}(c) \) follows from A1 and A2, using a codimension argument which we skip; see Peres [19] and Khoshnevisan [9, §4.7, p. 435] for details.

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