Experimental and Numerical Study of a Two-Stage Swirl Burner

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Abstract: In this study, we developed the design process and optimization of structural parameters of a new low-NOₓ burner based on low-NOₓ combustion technology and the flame stabilization principle. Firstly, on the basis of the two-stage swirl burner, we applied the fuel-graded combustion technology and introduced the central nozzle structure to explore the influence law of graded combustion on NOₓ emissions. Secondly, on the previously optimized structure, the matching law between the first- and second-stage cyclone blades is analyzed to obtain the optimum structural design solution for heat exchange efficiency and flame front length. Finally, a new conical blunt structure is introduced in conjunction with the flame stabilization principle, and we discuss the effects of different half cone angles on the flame stabilization, flame front length, and heat exchange efficiency of the burner. The research in this paper provides a reliable direction for the design optimization of low-NOₓ burners.

Keywords: low-NOₓ burner; flame stabilization; fuel-staging; optimal design

1. Introduction

Against the backdrop of energy-saving and emission reduction, the low-NOₓ swirl burner has become the main research direction in China [1]. When the burner is operating under rated conditions, the products after fuel combustion must meet the corresponding air pollutant emission standards based on environmental protection considerations [2]. For safety purposes, it must be ensured that backfiring or flame failure is not likely to occur. It can be seen that low pollutant emissions and flame stability must be achieved in the design or modification of the burner [3].

With the increasing interest in NOₓ control and the demand for guidelines for low-NOₓ burner designs, commonly used low-NOₓ combustion technology is divided into the following five categories: (1) graded combustion technology, characterized by mixing air and fuel in different areas in a certain proportion and completing combustion [4–6]; (2) low excess air coefficient combustion technology to make the combustion process as near as possible to the theoretical air volume [5,7]; (3) flue gas re-circulation technology, with the use of low temperature and low oxygen characteristics of flue gas, reducing the local high temperature in the furnace and resulting in the formation of a local reducing atmosphere, so as to inhibit the formation of NOₓ [8–10]; (4) lean premixed combustion technology, or the premixing of air and gas before combustion, which has the advantages of high temperature and high combustion intensity [11,12]; and (5) flameless combustion technology, a high-temperature, low-oxygen combustion technology, which is characterized by oxygen being heated to a certain temperature before reacting with the fuel, where the temperature difference in the furnace is small and the local high-temperature zone is reduced, thus reducing NOₓ emissions [13,14].

Many researchers have investigated low-NOₓ burner technology in practical applications using two main approaches: the use of a burner test facility or CFD modeling, which allows scholars to optimize multiple design parameters related to NOₓ emissions from...
industrial burners simultaneously. Zhao et al. [15] built an experimental platform to research the important role of fuel flow and equivalent volume ratio in generating non-linear thermoacoustic instabilities in swirl combustion chambers. Zhang et al. [16] investigated the effect of fuel mixture uniformity on the dynamic response of the flame. The study proved that the fuel/air mixture uniformity should be strictly controlled during the design of the combustion chamber to make the combustion chamber work safely and efficiently. Hadef et al. [17] used a Doppler anemometer (PDA) to compare the spatial distribution of droplets produced by the burner atomizing nozzles with two configurations of cyclonic flow: co-cyclonic and counter-cyclonic. The results show that the counter-cyclonic configuration atomizes better, produces high additional turbulence levels, and enhances fuel–air mixing efficiency. De Rosa’s experimental study of single-nozzle cyclonic combustion chamber structures identified that the overall flame outcome changed as the degree of flame wall interaction changed [18]. Legrand et al. [19] utilized stereo PIV measurements on a lean premix burner and observed that flame stability was better in the low cyclonic burner flow regime. Rosec’s experimental research into the field of high-oxygen fuels using thermal recirculation of exhaust gases in gas turbines has proved that this method significantly reduces CO, NO\textsubscript{x}, and soot emissions [20]. Yafei Zhang et al. [21] conducted combustion experiments with a real-time adjustable internal secondary air torch on a swirl burner, demonstrating that changing the torch angle can improve ignition and burnout and reduce NO\textsubscript{x} emissions under rich/lean combustion and load variations. Daood et al. [22] collected data by testing the main emissions in the flue gas and flame as well as the temperature and unburned carbon content of the ash throughout the chamber, confirming that oxygen-enriched combustion of the burner helped to control the products of NO\textsubscript{x} re-combustion.

In terms of numerical simulation, several authors have carried out computational fluid dynamics (CFD) simulations of the flow of combustion gas fuels in swirl burners. L.X. Zhou’s study of swirl combustion chamber combustion using large-eddy simulation (LES) found that the blunt-body structure changes the flow structure within the chamber, leading to a change in the flame structure of the gas combustion [23]. Wu et al. [24] carried out a study on the NO\textsubscript{x} emission characteristics of coal-fired boilers under staged combustion using Chemkin software. It was found that the conversion of fuel NO\textsubscript{x} at the exit of the main combustion zone increased with the stoichiometric ratio and was less influenced by the combustion temperature. Zhu et al. [25] performed numerical calculations based on RANS for a cyclonic combustion chamber and found that the central body jet improves the combustion performance.

Natural gas low-NO\textsubscript{x} burners are widely used because of their low NO\textsubscript{x} emissions, and swirl burners are used extensively in engineering combustion because of their reasonable airflow distribution with NO\textsubscript{x} generation, as well as their enhanced combustion and flame stabilization [26]. However, a review of the literature shows that there are few studies on the combustion performance of natural gas combustion chambers with two-stage blade swirlers, and it is of great value and significance to study the NO\textsubscript{x} emission flow field and characteristics of different combustors. In this paper, we examine the problems of poor flame stability, low heat exchange efficiency, and high NOx emissions of gas burners in the thermal system of asphalt mixing equipment, and we research the structural optimization design of natural gas burners. We used the NZG0B4000 as a study subject, a commercial combustor that uses natural gas as fuel and has a power range of 3500 KW to 28,000 KW. The influence of the two-stage blade structure on the performance of the combustion chamber is investigated through experiments and numerical simulations with the flame length, heat exchange efficiency, and NOx emission as evaluation indexes, and the structure of the burner is designed and optimized.

2. Experimental Setup and Methodology

The present experimental combustor consists of four major components, namely, a draught fan, gas pipe, swirler, and furnace, as shown in Figure 1. The mainstream air is
provided by the draught fan, and the gas pipe connecting to a gas pipeline transports fuel (natural gas). The furnace is a vaulted chamber with a transition cylinder on its inlet. It has a length of 11,290 mm. The thermocouples were installed at each test point to measure the temperature in the furnace. A cartesian coordinate is established and its origin point is located at the center of the combustor exit. The thermocouple is shown with T in Figure 1.

Figure 1. Schematic of the combustor geometry.

As shown in Figure 2a, the swirler comprises a double concentric swirler, consisting of a first-stage blade and a second-stage blade, with gas nozzles distributed evenly around the circumference of the second-stage blade. The air flowing through the swirler forms a rotating jet and the rest of the air forms a direct jet. Figure 2b shows the co-rotating swirler used in this experiment. The specific structural parameters of the combustor are shown in Table 1.

Table 1. Parameters for the combustor.

| Parameters                                      | Values |
|------------------------------------------------|--------|
| Diameter of combustor exit (mm)                | 830    |
| Blade number of first stage blade              | 28     |
| Blade number of secondary stage blade          | 30     |
| Install angle of first stage blade (°)          | 60     |
| Install angle of secondary stage blade (°)      | 65     |
| Vane thickness (mm)                            | 1      |

A swirler is usually characterized by the swirl number $S$, and the swirl number $S$ is used to reflect the rotating strength [27]. This dimensionless quantity is defined as the ratio
between the axial flux of angular momentum $G_\phi$ and the axial flux of axial momentum $G_x$ and is given by:

$$ S = \frac{G_\phi}{G_x R} $$  \hspace{1cm} (1)

where $R$ generally represents the jet outlet radius.

The chord length $l$ and the blade angle $\beta$ are constant for the axial swirler, so the axial velocity on the cross section can be approximately evenly distributed. Taking $R = R_1$, a simplified expression to calculate swirl number is given by:

$$ S = \frac{2}{3} \left[ 1 - \left( \frac{R_0}{R_1} \right)^3 \right]^{\frac{1}{2}} \tan \beta $$  \hspace{1cm} (2)

where $\beta$ is the blade angle between the normal direction of the blade and the direction of the airstream, and $R_0$ and $R_1$ are the inner and outer radii of the blade, respectively.

The whole experimental device is designed for both cold-flow and combustion experiments. According to the experimental specification requirements, the frequency of the draft fan is adjusted by the data measured by the Pitot tube (Product Model Number: JY-GD680) to achieve the same boundary conditions of the air inlet. The swirling vortex flowmeter (HP-LUX-DN100YF3E1B1) controls the methane flow rate, and the accuracy is 1% of the measurement range. Temperature distribution inside the combustor is measured by a 300 mm-long thermocouple (Platinum/rhodium alloy thermocouple, Type S Product Model Number: WRP-130). The measurement range is 0–1600 °C and the measurement accuracy is ±6 °C. The concentrations of emissions were obtained by the probe, i.e., flue gas analyzer (Testo-350). The sampling probe was mounted at the exit of the combustor. The measurement range is 0–500 ppm, the resolution is 0.1 ppm, and the accuracy is ±2 ppm. Other parameters used in the whole experiment are summarized in Table 2. The fuel type of the burner is natural gas, and its main component is methane. For ideal incompressible gas, the velocity inlet is adopted at the inlet, and the pressure outlet is adopted at the outlet. The non-slip boundary is adopted on the wall surface of the burner and the wall surface of the furnace.

**Table 2.** Parameters for the experiments.

| Parameters                              | Values  |
|-----------------------------------------|---------|
| Velocity of air inlet (m/s)             | 18.92   |
| Diameter of air inlet (m)               | 0.35    |
| Velocity of gas inlet (m/s)             | 28      |
| Diameter of gas inlet (m)               | 0.141   |
| Atmospheric temperature (K)             | 298     |
| Air inlet temperature (K)               | 298     |
| Excess air coefficient                  | 1.18    |
| Outlet pressure (Pa)                    | −70     |

The comparison between the temperature measured by the experiment and the results obtained by the numerical simulation is shown in Table 3. The numerical error between the two is less than 10%, and the simulation results have satisfactory reliability, so the selected values in this paper are verified. Figure 3 shows the flame shape of the burner in the furnace during the experiment.

**Table 3.** Comparison of experimental and numerical simulation results.

| Thermocouple Number | Experiment | Simulation | Error  |
|---------------------|------------|------------|--------|
| T1                  | 798 °C     | 827 °C     | 3.6%   |
| T2                  | 1092 °C    | 998 °C     | 8.6%   |
| T3                  | 985 °C     | 956 °C     | 2.9%   |
| T4                  | 967 °C     | 1008 °C    | 4.2%   |
3. Numerical Method

3.1. Realizable $k$-$\varepsilon$ Turbulence Model

The Realizable $k$-$\varepsilon$ turbulence model was chosen as the turbulent model [28]. The governing equation for the turbulent flow is given by:

$$\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_j}(\rho k u_j) = \frac{\partial}{\partial x_j} \left[ \left( \mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + G_k + G_b - \rho \varepsilon - Y_M + S_K$$

(3)

$$\frac{\partial}{\partial t}(\rho \varepsilon) + \frac{\partial}{\partial x_i}(\rho \varepsilon u_i) = \frac{\partial}{\partial x_j} \left[ \left( \mu + \frac{\mu_t}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right] + \rho C_1 S \varepsilon - \rho C_2 \frac{\varepsilon^2}{k + \sqrt{\varepsilon}} + C_{1\varepsilon} \frac{\varepsilon}{k} G_k + S_\varepsilon$$

(4)

where $C_1 = \max \left[ 0.43, \frac{\eta}{\eta + 5} \right]$, $\eta = S_k^2$, and $S = \sqrt{2 S_{ij} S_{ij}}$, where $G_k$ represents the generation of turbulence kinetic energy due to the mean velocity gradient. $G_b$ is the generation of turbulence kinetic energy due to buoyancy. $Y_M$ represents the contribution of the fluctuation incompressible turbulence to the overall dissipation rate. $C_2$ and $C_{1\varepsilon}$ are constant. $\sigma_k$ and $\sigma_\varepsilon$ are the turbulent Prandtl numbers for $k$ and $\varepsilon$, respectively. $S_K$ and $S_\varepsilon$ are user-defined source terms.

3.2. Radiative Transfer Model

The radiative transfer equation [29] for an absorbing, emitting, and scattering medium at position in $\vec{r}$ the $s$ direction is:

$$\frac{dI(\vec{r}, s)}{ds} + (a + \sigma_s) I(\vec{r}, s) = an^2 \sigma T^4 \int_0^{4\pi} I(\vec{r}, s') \phi(s, s') d\Omega'$$

(5)

where $\vec{r}$ is position vector, $\vec{s}$ is direction vector, $\vec{s'}$ is scattering direction vector, $s$ is path length, $a$ is absorption coefficient, $n$ is refractive index, $\sigma_s$ is scattering coefficient, $\sigma$ is Stefan–Boltzmann constant radiation intensity, $T$ is local temperature, $\phi$ is phase function, $\Omega'$ is solid angle, and $(a + \sigma_s)$ is the optical thickness or opacity of the medium. The refractive index $n$ is important when considering radiation in semi-transparent media.

3.3. NO$_x$ Formation Model

In laminar flames and at the molecular level within turbulent flames, the formation of NO$_x$ can be attributed to four distinct chemical kinetic processes: thermal NO$_x$ formation, prompt NO$_x$ formation, fuel NO$_x$ formation, and intermediate N$_2$O [30]. The present numerical simulation mainly considers thermal NO$_x$ and fast NO$_x$. 

Figure 3. The flame of the burner in the furnace during the experiment.
The formation of thermal NO\textsubscript{x} is determined by a set of highly temperature-dependent chemical reactions known as the extended Zeldovich mechanism. The principal reactions governing the formation of thermal NO\textsubscript{x} from molecular nitrogen are as follows:

\[ O + N_2 \rightarrow N + NO \]  
\[ N + O_2 \rightarrow O + NO \]  

Prompt NO\textsubscript{x} formation is most prevalent in rich flames. The actual formation involves a complex series of reactions and many possible intermediate species. The route now accepted is as follows:

\[ CH + N_2 \rightarrow HCN + N \]  
\[ N_2 + O \rightarrow NO + O \]  
\[ HCN + OH \rightarrow CN + H_2O \]  
\[ CN + O_2 \rightarrow NO + CO \]

### 3.4. Combustion Reaction Model

Fuel combustion is influenced by the interaction between turbulent flow and chemical reactions, with turbulence intensifying the mixing and influencing the time-averaged chemical reaction rate, and conversely, the exothermic process of combustion influencing turbulent flow. This paper adopts the more widely used eddy dissipation model (EDM). The reaction rate is calculated from the following equation, which takes the minimum of the two:

\[
R_{i,r} = \nu_{i,r}^' M_r A \rho \varepsilon \min \left[ \frac{\omega_R}{\nu_{i,r}^' M_r}, \frac{\sum \omega_P}{\nu_{i,r}^' M_r} \right]
\]  

where the reaction rate of component i in chemical reaction r is \( R_{i,r} \), and \( \nu_{i,r}^' \) is the coefficient of chemical equivalence of reactants and products of component i in reaction r. \( M_r \) is the relative molecular mass of the component, \( N \) is the total number of products, and \( \omega \) is the mass fraction of components.

### 3.5. Numerical Meshes

The fluid area of the combustor has a complex structure, so the fluid area is divided. The unstructured tetrahedral mesh is used to divide the unimportant fluid area, such as the non-combustion area of the furnace chamber. The overall mesh distortion is less than 0.8, the mesh quality is greater than 0.3, and the total number of mesh cells is 9.86 million. The grid is shown in Figure 4.

![Figure 4. Overall grid diagram of burner numerical simulation area.](image)

### 3.6. Numerical Solution Methods

Two solution methods are provided in Fluent software, the pressure-based solver and the density-based solver, either of which must be used to solve the fundamental equations. The grid partitioning discretizes the fluid region into different control bodies and integrates
over each body to create the algebraic equations from which the fundamental unknown variables are solved until the calculation converges. In this paper, SIMPLE is used to solve for the coupling of velocity and pressure, and the QUICK format is used to discretize the momentum equations.

Due to a large number of grids in this model, the other equations are calculated in first-order windward format and then in the second-order windward format to speed up the convergence of the calculations. The difference between the two is the expansion of the Taylor formula, which has higher accuracy than the first-order windward format.

4. Simulation and Analysis

4.1. The Effect of a Gas Classification Strategy with a Central Nozzle on the Flow Field and Combustion Characteristics

Figure 5a shows a two-stage cyclone structure, with the angles of the first stage and second stage blades being 65° and 20°, respectively. This section of the paper is based on the structure of the D65-30, using gas classification combustion technology to carry out structural solutions. The improved cyclone structure is shown in Figure 5 and is named CD65-30. Compared to D65-30, the latter has additional structures such as central gas piping and gas nozzles. The total number of nozzles is 12, with six central gas nozzles and six peripheral gas nozzles each. The purpose of this improved construction scheme is to increase the efficiency of the heat exchange between the flame and the aggregate by encouraging more gas and air to be burned in the upstream area of the furnace. At the same time, the fuel distribution is carried out, i.e., the gas to air equivalent ratio in the primary air area is controlled and the combustion in the primary blade area is set to thin combustion of the fuel using the gas classification combustion technique. This is verified by simulating the flow of the burner in CFD, with axial and radial velocity distributions at different positions of the burner.

Figure 5. Three-dimensional structure of (a) D65-30 and (b) CD65-30.

It can be seen from Figure 6a that there is a recirculation zone near the end of the central gas pipe for CD65-30, which is formed by the high-speed air flowing through the gap between the swirler and the central gas pipe, and the range of this area is very small, shown in Figure 6b. Based on Figure 6c,d, it can be seen that the recirculation zone formed on the periphery of CD65-30 is larger than D65-30, so the flame stability is improved. In addition, the axial velocity of D65-30 is greater than that of CD65-30, and higher axial velocity is detrimental to shortening the length of the flame front.

From Figure 7a, it can be seen that the tangential velocity of CD65-30 increases, then decreases and then increases again along the radial direction at this section (x = −0.55 m). The peak of the tangential velocity occurs in the region where the radial distance is less than 0.05 m. This area is the reflux zone near the central gas pipe. The effect of the central recirculation zone makes the overall tangential velocity of CD65-30 less than that of D65-30. From Figure 7b–d, it can be seen that the secondary wind region produces negative tangential velocity, that is, the direction of rotation of the airflow here is opposite to the
direction of rotation of the main airflow. The reason is that the angles of the primary blade and the secondary blade are different after the cyclone blade, and there is a tangential velocity difference.

**Figure 6.** Axial velocity of D65-30 and CD65-30 at different axial positions.

**Figure 7.** Tangential velocity of D65-30 and CD65-30 at different axial positions.
Figure 8 shows that the nozzles are arranged in such a way that the gas and air are further mixed for CD65-30, and the boundary layer between the two regions undergoes intense mass and energy exchange. The energy produced by the combustion in the central pipe area is gradually transferred to the periphery, which is beneficial to the stability of the flame. Compared with the length of the flame, it is found that the difference between the two is very small. In contrast to the D65-30, the gas injected by the central nozzle burns near the downstream of the cyclone and forms a narrow, high-temperature zone. The arrangement of the CD65-30 gas nozzles in the burner’s cowl allows further graded mixing of gas and air, with intense mass and energy exchange taking place on the boundary layer between the two zones under strongly turbulent combustion. For the CD65-30, the high-temperature flue gas produced by combustion in the central zone, the energy contained in the flue gas is transferred to the periphery step by step, which is advantageous for flame stabilization. The heat produced by the D65-30 after combustion in the peripheral secondary air zone is transferred more to the peripheral medium by the strong turbulence, and only a small amount of heat energy is diffused to the interior.

![Figure 8. Temperature distribution of D65-30 and CD65-30 on the center cross section.](image)

Comparing the temperature near the wall of the furnace from Figure 9, the temperature of D65-30 is slightly higher than that of CD65-30, which means that the setting of the nozzle does not significantly improve the heat exchange efficiency between the flame and the aggregate. The position near the wall is also where the thermocouple is installed on the upper wall of the furnace chamber.

![Figure 9. Temperature distribution of D65-30 and CD65-30 near the furnace wall.](image)

As can be seen from Figure 10, the D65-30 has an overall higher NO\textsubscript{x} concentration in the furnace chamber than the CD60-30. A large area with a NO\textsubscript{x} concentration above 70 ppm
exists downstream of the D65-30 burner. A comparison of the temperature profiles in Figure 8 shows that the location of this area of high NO\textsubscript{x} concentration is downstream of the high-temperature zone in the furnace chamber where the local temperature exceeds 2000 K. Due to the local high temperatures in this area, the nitrogen in the air is rapidly converted to thermodynamic NO\textsubscript{x}. CD65-30 also has an area of high local NO\textsubscript{x} concentration in the same area, but the area and value of this area are smaller in comparison to D65-30.

![Figure 10. Local distribution of NO\textsubscript{x} concentration (ppm) of D65-30 and CD65-30 on the central cross section.](image)

As shown in Table 4, the NO\textsubscript{x} concentration of the CD65-30 is lower than that of the D65-30 at the exit of the furnace due to the graded combustion of the gas, the introduction of a central gas nozzle, and the multi-layer arrangement of the nozzles, which allows for further thinning of the fuel, thus reducing the NO\textsubscript{x} concentration.

| Name     | NO\textsubscript{x} Concentration (ppm) |
|----------|----------------------------------------|
| D65-30   | 70.8                                   |
| CD65-30  | 57.3                                   |

In summary, the arrangement of the central gas pipe and the nozzle has no obvious effect on the improvement of the flame front length, but the introduction of gas-staged combustion technology has optimized the NO\textsubscript{x} emissions, so it is desirable to improve the burner in terms of NO\textsubscript{x} emissions.

4.2. The Influence of the Angle of the Two-Stage Swirl Blade on the Flow and Combustion Characteristics

Scholars have pointed out that the dual-stage swirler needs to be reasonably matched, otherwise it will not make many contributions to the combustion process [31]. Thus, we conducted a parametric analysis of the second-stage swirl blade angle to explore the matching rule with the first-stage blade angle of based on the burner (CD65-30). A three-dimensional diagram of different secondary-stage swirl blade angles is shown in Figure 11.

After the statistics of the airflow rates through the swirler, it was found that as the angle decreased, the secondary air volume gradually increased, and the tertiary air volume remained unchanged. It can be seen from Figure 12a that there is a recirculation zone around the central axis, and the reflow velocity gradually decreases as the secondary-stage blade angle decreases. In addition, the position of the peak axial velocity gradually shifts outward. It can be seen from Figure 12b that in the vicinity of the radial position of \( r = 0.25 \) m, the peak axial velocity of CD65-30 is the smallest. However, the viscous
force causes the axial velocity to decrease in comparison to the other two burners in this region of the axial velocity peak. In addition, there is a recirculation zone on the periphery and its corresponding position is near the peripheral gas pipe, and the reflow velocity of CD55-30 in this recirculation zone is the largest. Based on Figure 12c,d, the recirculation zone formed by CD55-30 and CD45-30 in the periphery disappears, which means that the larger the angle of the secondary-stage swirl blades, the more favorable the formation of the recirculation zone.

![Figure 11. Three-dimensional diagram of different secondary-stage swirl blade angles.](image1)

![Figure 12. Axial velocity of CD65-30, CD55-30, and CD45-30 at different axial positions.](image2)

It can be seen from Figure 13a that the tangential velocity first increases and then decreases when the radial distance is less than 0.05 m. Since the potential flow rotation zone and the quasi-rigid body rotation zone are formed in this recirculation zone, the tangential velocity distribution is opposite in these two zones. In addition, the negative values appear for tangential velocity, indicating that the airflow is not symmetrical about the central axis. We can see from Figure 13a–d that the overall tangential velocity gradually decreases as the angle decreases, and the effect of the angle on the tangential velocity of the airflow dominates despite the difference in air volume at all stages. Therefore, the greater the tangential velocity of the swirling flow, the more uniform the mixing is.
Referring to the temperature contour of 1800 K, the location of the temperature contour at 1800 K is taken as the location of the flame front. It is found from Figure 14 that the length of the flame front gradually shortens as the angle decreases. The presence of high temperature in the hood is beneficial to the ignition. Considering the effect of combustion expansion on the flame, CD55-30 and CD45-30 are easily blown out. From the perspective of flame stability, the blade’s angle of the two does not match the first-stage blades. The locally uneven high temperature of CD55-30 and CD45-30 promotes the formation of NOx when the air excess coefficient is greater than 1.
It can be seen from Figure 15 that the temperature of CD45-30 near the front wall of the furnace is higher than the other two. The axial velocity of the CD45-30 is the smallest from Figure 12, so the mixing and combustion position is closer to the burner outlet. In addition, when the axial distance is less than 2.5 m, the temperature of CD65-30 is higher than that of CD55-30, which means that the heat exchange efficiency is higher than that of CD55-30.

![Temperature distribution of CD65-30, CD55-30, and CD45-30 near the furnace wall.](image)

**Figure 15.** Temperature distribution of CD65-30, CD55-30, and CD45-30 near the furnace wall.

As can be seen in Figure 16, as the angle decreases, the overall NO\(_x\) concentration gradually increases. In addition, NO\(_x\) is generated for CD55-30 and CD45-30 in the hood, because the amount of secondary air increases as the angle decreases to form a rapid NO\(_x\).

![NO\(_x\) local distribution of CD65-30, CD55-30, and CD45-30 on the central cross section.](image)

**Figure 16.** NO\(_x\) local distribution of CD65-30, CD55-30, and CD45-30 on the central cross section.

From Table 5, it is found that NO\(_x\) emissions have improved for CD65-30, but the impact of the induced draft fan on the flow field and combustion characteristics has not been effectively improved. In addition, the recirculation zone formed near the peripheral nozzle is weakened, so the flame stability is reduced.
Table 5. Average NOx concentration at the furnace exit for CD65-30, CD55-30, and CD45-30.

| Name   | NOx Concentration (ppm) |
|--------|-------------------------|
| CD65-30| 57.3                    |
| CD55-30| 76.2                    |
| CD45-30| 84.6                    |

4.3. The Influence of Bluff Body on Flow and Combustion Characteristics

From the above problems, we propose a bluff body structure combined with the CD65-30 to solve the above problems in a comprehensive and balanced manner, as shown in Figure 17. A cone-shaped bluff body is arranged on the periphery of the swirler, and the bluff body will force the flow direction of the fluid to change.

![Figure 17. Three-dimensional structure of swirler with bluff structure.](image)

We parametrically studied the semi-cone angle of the bluff body with intervals of 10°, and named D65-C45, D65-C35, and D65-C25, respectively. We found that the resistance of the bluff body to the flow of the tertiary air decreases as the angle decreases. However, the increased airflow rates flowing through the swirler weakens the effect of the induced fan on the downstream fluid distribution, intensifying the mixing of gas and air and thus improving the uniformity.

From Figure 18a–d, it is found that the axial velocity of the D65-C45 as a whole is greater than that of the other two burners. The primary airflow is greater than the secondary airflow at different angles, thus showing that the primary airflow dominates the movement of the total airflow. Moreover, the axial velocity of the D65-C45 is greater than that of the other two burners at different angles. Combining Figure 18b–d, it can be seen that the recirculation zone of D65-C25 is the largest in the vicinity of the bluff body.

In addition, a comparison of Figure 12c,d and Figure 18c,d reveals that the area of the recirculation zone formed by D65-C25 is larger than that of CD65-30, so this increases the energy of the recirculation zone and accelerates the reaction rate of the mixed gas at the boundary, which means that the proposed conical bluff body structure improves the flame stability.

It can be seen from Figure 19 that D65-C45 has the largest tangential velocity at any point, which is consistent with the relevant conclusions above. As more air flow passes through the swirler, this will generate a rotary motion. According to the conservation of momentum, with greater rotational angle kinetic energy, the air can quickly spread out. The tangential velocity affects the mixing in diffusion combustion, which in turn affects the temperature distribution and NOx generation throughout the combustion process.
Figure 18. Axial velocity distribution of D65-C45, D65-C35, and D65-C25 at different axial positions.

Figure 19. Tangential velocity distribution of D65-C45, D65-C35, and D65-C25 at different axial positions.
As can be seen from Figure 20, the shortest flame front length is found for D65-C45 using the 1800 K contour as a reference. As the half cone angle increases, the flame front length increases and then decreases. The tangential velocity of the D65-C45 is highest at the exit of the chamber, and the mixing and combustion of the gas and air occur on both sides of the central axis due to the high centrifugal forces. This is reflected in the fact that the high-temperature area of the chamber is not concentrated on the central axis, which is favorable for the heat exchange between the flame and the aggregate. In addition, the D65-C45 has a larger range of high-temperature areas present in the header compared to the CD65-30. It can be deduced from this that the conical blunt body structure strengthens the ability to roll up the high-temperature flue gases in the return zone near the gas nozzle, resulting in increased energy in the area, which plays an important role in the stabilization of the flame root.

![Figure 20. Temperature distribution of D65-C45, D65-C35, and D65-C25 in the center section.](image)

It can be seen that the temperatures for D65-C45 are highest near the wall and about 100 K above CD65-30 from Figure 21, indicating heat exchange between the flame and aggregate. The heat exchange efficiency between them is improved by the proposed blunt structure. Therefore, setting up a bluff body improves the airflow distribution in the front of the swirler, which is of great practical importance.

![Figure 21. Temperature distribution of D65-C45, D65-C35, and D65-C25 near the furnace wall.](image)
It can be seen from Figure 22 that as the half-cone angle of the tapered bluff body decreases, the NO\textsubscript{x} concentration changes in a decreasing trend, and the NO\textsubscript{x} concentration of D65-C45 is the highest. When the half-cone angle of the cone-shaped bluff body is 45\degree, the formed recirculation zone is the largest, and the energy is high after a large amount of smoke is entrained.

Figure 22. NO\textsubscript{x} local distribution of D65-C45, D65-35, and D65-C25 on the central cross section.

At the same time, D65-45 has the most energy in the recirculation zone near the bluff body compared to the other two burners, so it can ignite a larger amount of mixed gas, which speeds up the combustion process and quickly releases heat. It can be seen from Figure 22 that the temperature of D65-C45 in the combustor headcover is higher than the other two. In addition, due to the larger tangential velocity, the surface area of the flame front during combustion is increased. When the excess air coefficient is greater than 1, the excess oxygen quickly generates thermal NO\textsubscript{x} at high temperatures. Therefore, the NO\textsubscript{x} of D65-C45 has the largest amount of production. For D65-C35 and D65-C25, since the tangential velocity of the total airflow is greatly reduced compared with D65-C45, the weaker airflow rotation ability weakens the gas mixing effect and slows down the combustion process. The main area is concentrated on the central axis, so the generation of NO\textsubscript{x} is also mainly near the central axis.

In comparison with Tables 4 and 6, the presence of the blunt body structure increases the NO\textsubscript{x} emissions, but the flame stability, flame front length, and temperature distribution near the upstream wall of the chamber show that the D65-C45 offers the best overall performance compared to the D65-30 and CD65-30 when the conical blunt structure solution is improved.
Table 6. Average NO\textsubscript{x} concentration at the furnace exit for D65-C45, D65-C35, and D65-C25.

| Name     | NO\textsubscript{x} Concentration (ppm) |
|----------|----------------------------------------|
| D65-C45  | 64.9                                   |
| D65-C35  | 58.4                                   |
| D65-C25  | 57.9                                   |

5. Conclusions

Based on the NO\textsubscript{x} strategy for staged burners, we investigated the optimal design of natural gas burners in combination with computational fluid dynamics analysis and experiments. The following conclusions were mainly obtained:

(1) By setting a central gas nozzle on the basis of a two-stage swirl structure, the NO\textsubscript{x} concentration is reduced from 70.8 ppm to 57.3 ppm with essentially constant airflow at all stages due to the gas-staged combustion technology. However, the length of the flame front, flame stability, and heat exchange efficiency do not change significantly.

(2) Parametric analysis of the two-stage swirler blade shows that, with the first stage blade angle unchanged, a decrease in the angle of the second stage blade reduces the tangential velocity of the total airflow and weakens the gas mixing effect of the two-stage swirler. As the angle of the second stage blade decreases, the flame front length gradually shortens, flame stability gradually decreases, and the concentration of NO\textsubscript{x} gradually increases. The decrease in the angle of the second stage blade gives the heat exchange efficiency a trend of decreasing first and then increasing.

(3) To stabilize the flame and reduce the flame front length, a new blunt body structure, D65-C45, is proposed, which reduces the NO\textsubscript{x} concentration at the furnace exit to 64.9 ppm, corresponding to a flame front length of approximately 4.8 m.

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