Rieffel deformation and twisted crossed products

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Abstract

To a continuous action of a vector group on a C*-algebra, twisted by the imaginary exponential of a symplectic form, one associates a Rieffel deformed algebra as well as a twisted crossed product. We show that the second one is isomorphic to the tensor product of the first one with the C*-algebra of compact operators in a separable Hilbert space and we indicate some applications.

1 Introduction

In order to provide a unified framework for a large class of examples in deformation quantization, Marc Rieffel [15] significantly extended the basic part of the Weyl pseudodifferential calculus. Rieffel’s calculus starts from the action Θ of a finite-dimensional vector space Ξ on a C*-algebra A, together with a skew-symmetric linear operator J : Ξ → Ξ that serves to twist the product on A. Using J one defines first a new composition law # on the set of smooth elements of A under the action and then a completion is taken in a suitable C*-norm. The outcome is a new C*-algebra A, also endowed with an action of the vector space Ξ. The corresponding subspaces of smooth vectors under the two actions, A∞ and A∞, respectively, coincide. In [15] the functorial properties of the correspondence A → A are studied in detail and many examples are given.
It is also shown that one gets a strict deformation quantization of a natural
Poisson structure defined on $\mathcal{A}$ by the couple $(\Theta, J)$. Assuming $J$ non-degenerate (so it defines a symplectic form on $\mathcal{E}$), one gets a twisted action $(\Theta, \kappa)$ of $\mathcal{E}$ on the $C^*$-algebra $\mathcal{A}$, where $\kappa$ is the 2-cocycle on $\mathcal{E}$ given by $(X, Y) \mapsto \kappa(X, Y) := \exp(iX \cdot J Y)$. To such a twisted action one associates canonically \cite{12, 13} a twisted crossed product $C^*$-algebra $\mathcal{A} \rtimes_{\Theta} \mathcal{E}$, whose representations are determined by covariant representations of the quadruplet $(\mathcal{A}, \Theta, \kappa, \mathcal{E})$.

In the present article we are going to show that the two $C^*$-algebras $\mathcal{A}$ and $\mathcal{A} \rtimes_{\Theta} \mathcal{E}$ that can be constructed from the data $(\mathcal{A}, \Theta, \kappa, \mathcal{E})$ are actually stably isomorphic. This happens in a particularly precise way: one has an isomorphism (called the canonical mapping) $M : \mathcal{H} \otimes \mathcal{A} \rightarrow \mathcal{A} \rtimes_{\Theta} \mathcal{E}$, where $\mathcal{H}$ is an elementary $C^*$-algebra, i.e. it is faithfully represented as the ideal of all compact operators in a separable Hilbert space. The mapping $M$ is naturally defined first between convenient Fréchet subalgebras (vector-valued Schwartz spaces); the extension to a $C^*$-isomorphism needs a non-trivial isometry argument.

Such stable isomorphism has standard consequences \cite{14}: the (closed, bi-
sided self-adjoint) ideals of the two algebras $\mathcal{A}$ and $\mathcal{A} \rtimes_{\Theta} \mathcal{E}$ are in one-to-one correspondence, the spaces of primitive ideals are homeomorphic and the two representation theories are identical. By using basic information about the twisted crossed product, we also get a simple proof of the known fact \cite{17, 8} that the $K$-groups of the Rieffel deformed algebra $\mathfrak{A}$ are the same as those of the initial algebra $\mathcal{A}$. A covariant morphism $\mathcal{R} : (\mathcal{A}^1, \Theta^1) \rightarrow (\mathcal{A}^2, \Theta^2)$ can be raised both to a morphism $\mathfrak{A} : \mathfrak{A}^1 \rightarrow \mathfrak{A}^2$ and to a morphism $\mathcal{R} \times : A^1 \rtimes_{\Theta^1} \mathcal{E} \rightarrow A^2 \rtimes_{\Theta^2} \mathcal{E}$. The canonical mappings $M^1, M^2$ have the intertwining property $\mathcal{R} \times \circ M^1 = M^2 \circ (\text{id} \otimes \mathfrak{A})$.

When the initial algebra $\mathcal{A}$ is commutative, it is associated by Gelfand theory with a locally compact topological dynamical system $(\Sigma, \Theta, \mathcal{E})$. Under some assumptions on this system, one can get information on the primitive ideal space of the $C^*$-algebra $\mathfrak{A}$. A choice of an invariant measure on $\Sigma$ leads to $L^2$-orthogonality relations for the canonical mapping $M$. We hope to continue to investigate the canonical mappings in the commutative case (the one closest in spirit with traditional pseudodifferential theory), having in view a more detailed study of representations, modulation spaces and applications to spectral analysis \cite{9}.

## 2 Involutive algebras associated to a twisted $C^*$-dynamical system

We shall recall briefly, in a slightly particular setting, some constructions and results concerning twisted crossed products algebras and Rieffel’s pseudodifferential calculus.

The common starting point is a $2n$-dimensional real vector space $\mathcal{E}$ endowed with a symplectic form $[\cdot, \cdot]$. When needed we are going to suppose that $\mathcal{E} =$
suitably defined by oscillatory integral techniques. Thus one gets a
action Θ leaves

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one associates canonically a (Hilbert module techniques; we are going to call

the product

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To quantize the above structure, one keeps the involution but introduce on

where

This is the enveloping

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-algebra of

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quadruplet

(2.1)

k : Ξ × Ξ → ℤ := \{λ ∈ ℤ | |λ| = 1\}, \quad \kappa(X, Y) := \exp\left(-\frac{i}{2} [X, Y]\right)

and notice that it is a group 2-cocycle, i.e. for all \(X, Y, Z \in Ξ\) one has

\[
\kappa(X, Y) \kappa(X + Y, Z) = \kappa(Y, Z) \kappa(X, Y + Z), \quad \kappa(X, 0) = 1 = \kappa(0, X).
\]

Thus the classical data is converted into \((A, Θ, Ξ, \kappa)\), a very particular case of
twisted \(C^\ast\)-dynamical system [12] [13]. To any twisted \(C^\ast\)-dynamical system
one associates canonically a \(C^\ast\)-algebra \(A ∗_θ Ξ\) (called twisted crossed product).
This is the enveloping \(C^\ast\)-algebra of the Banach \(\ast\)-algebra \((L^1(Ξ; A), ⊞, \| \cdot \|_1)\),
where

\[
\| G \|_1 := \int_Ξ dX \| G(X) \|_A, \quad G^0(X) := G(-X)^\ast
\]

and (symmetrized version of the standard form, cf. Remark [13]

\[
(G_1 ∘ G_2)(X) := \int_Ξ dY \kappa(X, Y) Θ_\lambda (Y - X)/2 \left[ G_1(Y) \right] Θ_Y/2 \left[ G_2(X - Y) \right]. \quad (2.2)
\]

In [12] [13]. \(A\) is supposed separable; since our cocycle is explicit and very simple,
this will not be needed here.

To quantize the above structure, one keeps the involution but introduce on \(A^∞\)
the product

\[
f \# g := 2^{2n} \int_Ξ \int dY dZ \ e^{2 i [Y, Z]} Θ_Y(f) Θ_Z(g), \quad (2.4)
\]

suitably defined by oscillatory integral techniques. Thus one gets a \(\ast\)-algebra
\((A^∞, \#, \ast)\), which admits a \(C^\ast\)-completion \(\mathfrak{A}\) in a \(C^\ast\)-norm \(\| \cdot \|_\mathfrak{A}\) defined by
Hilbert module techniques; we are going to call \(\mathfrak{A}\) the \(R\)-deformation of \(A\). The
action Θ leaves \(A^∞\) invariant and extends to a strongly continuous action on
the $C^*$-algebra $A$, that will also be denoted by $\Theta$. The space $A^\infty$ of $C^\infty$-vectors coincide with $A^\infty$, even topologically, i.e. the family $\{\mathcal{A}\}$ on $A^\infty = A^\infty$ is equivalent to the family of semi-norms

$$|f|_k^A := \sum_{|\alpha|=k} \frac{1}{\alpha!} \| \partial^\alpha_N [\Theta X (f)]_{X=0} \|_A = \sum_{|\alpha|=k} \frac{1}{\alpha!} \| \delta^\alpha(f) \|_A, \quad k \in \mathbb{N}. \quad (2.5)$$

An important particular case is obtained when $A$ is the $C^*$-algebra $BC_u(\Xi)$ of bounded uniformly continuous functions on the group $\Xi$, which is invariant under translations, i.e. if $a \in A$ and $X \in \Xi$, then $[T_X(a)](\cdot) := a(\cdot - X) \in A$. Notice that the $*$-algebra of smooth vectors coincides with $BC^\infty(\Xi)$, the space of all smooth complex functions on $\Xi$ with bounded derivatives of every order. In this case Rieffel's construction, done for $\Theta = T$, reproduces essentially the standard Weyl calculus; we are going to use the special notations $\sharp$ (instead of #) for the corresponding composition law and $\mathcal{B}(\Xi)$ for the R-deformation of $BC_u(\Xi)$.

One can also consider $C^*$-subalgebras $A$ of $BC_u(\Xi)$ that are invariant under translations. An important one is $C_0(\Xi)$, formed of all the complex continuous functions on $\Xi$ that decay at infinity. Its Rieffel deformation will be denoted by $\mathcal{K}(\Xi)$; it contains the Schwartz space $\mathcal{S}(\Xi)$ densely. By Example 10.1 and Proposition 5.2 in [15] it is elementary, i.e isomorphic to the $C^*$-algebra of all compact operators in a separable Hilbert space.

Following [15], we introduce the Fréchet space $\mathcal{S}(\Xi; A^\infty)$ composed of smooth functions $F : \Xi \to A^\infty = A^\infty$ with derivatives that decay rapidly with respect to all $| \cdot |_k^A$. The relevant seminorms on the space $\mathcal{S}(\Xi; A^\infty)$ are $\{\| \cdot \|^{k,\beta,N}_{A,\Xi} \}$ where

$$\| F \|^{|k,\beta,N}_{A,\Xi} := \sup_{X \in \Xi} \{(1 + |X|)^N \| (\partial^\beta F)(X) \|_A\}, \quad (2.6)$$

and the index $A$ can be replaced by $A$, by the argument above. We are going to use repeatedly the identification of $\mathcal{S}(\Xi; A^\infty)$ with the topological tensor product $\mathcal{S}(\Xi) \hat{\otimes} A^\infty$ (recall that the Fréchet space $\mathcal{S}(\Xi)$ is nuclear). On it (and on many other larger spaces) one can define obvious actions $\Xi := T \otimes 1$ and $\Xi \otimes \Theta$ of the vector spaces $\Xi$ and $\Xi \times \Xi$, respectively. Explicitly, for all $A,Y,X \in \Xi$, one sets $[\Xi_A(F)](X) := F(X - A)$ and $[(\Theta_A \otimes Y) F](X) := Y [F(X - A)]$. Then on $\mathcal{S}(\Xi; A^\infty)$ one can introduce the composition law

$$(F_1 \square F_2)(X) = 2^{2n} \int_{\Xi} \int_{\Xi} dAdB e^{-2i[A,B]} [\Xi_A(F_1)](X) \# [\Xi_B(F_2)](X) = \quad (2.7)$$

and

$$= 2^{4n} \int_{\Xi} \int_{\Xi} \int_{\Xi} dA dB dY dZ e^{-2i[A,B]} e^{2i[Y,Z]}$$

$$[(\Theta_A \otimes \Theta_Y) (F_1)](X) [(\Theta_B \otimes \Theta_Z) (F_2)](X). \quad (2.8)$$

Notice that the last expression should be interpreted as an oscillatory integral [15] and that it involves the multiplication in the $C^*$-algebra $A$. If the involution is given by $F^\Xi(X) := F(X)^*$, $\forall X \in \Xi$, it can be shown that one gets a Fréchet $^*$-algebra.
Remark 2.1. We recall that $\mathcal{A}^\infty = \mathfrak{A}^\infty$, even topologically, but the algebraic structures are different. When the forthcoming arguments will involve the composition $\#$, in order to be more suggestive, we will use the notation $\mathcal{F}(\Xi; \mathfrak{A}^\infty)$. In other situations the notation $\mathcal{F}(\Xi; \mathcal{A}^\infty)$ will be more natural. For instance, it is easy to check that $\mathcal{F}(\Xi; \mathcal{A}^\infty)$ is a (dense) $^*$-subalgebra of the Banach $^*$-algebra $(L^1(\Xi; \mathcal{A}), \circ, \| \cdot \|_1)$, which is defined in terms of the product $\cdot$ on $\mathcal{A}$ and has a priori nothing to do with the composition law $\#$. Proposition 4.2 is a good illustration for this distinction.

Remark 2.2. One can also consider $BC_u(\Xi; \mathfrak{A})$, the $C^*$-algebra of all bounded and uniformly continuous functions $F: \Xi \to \mathfrak{A}$. Rieffel deformation can also be applied to the new classical data $(BC_u(\Xi; \mathfrak{A}), \Xi, \Xi, [- , -])$, getting essentially (2.7) as the corresponding composition law. By using the second part (2.8) of the formula, this can also be regarded as the Rieffel composition constructed from the extended twisted $C^*$-dynamical system $(BC_u(\Xi; \mathfrak{A}), \mathcal{T} \otimes \Theta, \Xi \times \Xi, \pi \otimes \kappa)$. This will not be needed in this form. But we are going to use below the fact that for elements $f, g \in \mathfrak{A}^\infty$, $a, b \in \mathcal{F}(\Xi)$ one has $(a \otimes f) \square (b \otimes g) = (b \sharp a) \otimes (f \# g)$, so $\square$ can be seen as the tensor product between $\#$ and the law opposite to $\sharp$. By Proposition 2.1 in [16], one can identify $\mathcal{F}(\Xi) \otimes \mathfrak{A}$ with the $\mathcal{R}$-deformation of $C_0(\Xi) \otimes \mathfrak{A} \equiv C_0(\Xi; \mathfrak{A})$ and $\mathcal{B}(\Xi) \otimes \mathfrak{A}$ with the $\mathcal{R}$-deformation of $BC_u(\Xi) \otimes \mathfrak{A}$.

3 The Schrödinger representation

We are going to denote by $\mathcal{B}(\mathcal{M}, \mathcal{N})$ the space of all linear continuous operators acting between the topological vector spaces $\mathcal{M}$ and $\mathcal{N}$ and use the abbreviation $\mathcal{B}(\mathcal{M})$ for $\mathcal{B}(\mathcal{M}, \mathcal{M})$.

Let us recall that $\mathfrak{X}$ is a finite-dimensional vector space. The corresponding Heisenberg algebra $\mathfrak{h}_\mathfrak{X} = \mathfrak{X} \times \mathfrak{X}^* \times \mathbb{R}$ is the Lie algebra with the bracket

$$[(x, \xi, t), (y, \eta, s)] := (0, 0, y \cdot \xi - x \cdot \eta).$$

We use notations as $\tilde{X} = (x, \xi, t)$ and $X = (x, \xi)$. The Heisenberg group $\mathbb{H}_\mathfrak{X}$ is just $\mathfrak{h}_\mathfrak{X}$ thought of as a group with the multiplication $*$ defined by

$$\tilde{X} * \tilde{Y} = \tilde{X} + \tilde{Y} + \frac{1}{2} [\tilde{X}, \tilde{Y}], \quad \tilde{X}, \tilde{Y} \in \mathbb{H}_\mathfrak{X}.$$

The unit element is $0 \in \mathbb{H}_\mathfrak{X}$ and the inversion mapping given by $\tilde{X}^{-1} := -\tilde{X}$.

The Schrödinger representation is the unitary representation $\Pi: \mathbb{H}_\mathfrak{X} \to \mathcal{B}(\mathcal{L})$ in the Hilbert space $\mathcal{L} := L^2(\mathfrak{X})$, defined by

$$[\Pi(\tilde{X})u](y) = [\Pi(x, \xi, t)u](y) = e^{i(y \cdot \xi + \frac{1}{2}x \cdot \xi + t)} u(y + x) \quad \text{for a.e. } y \in \mathfrak{X} \quad (3.1)$$

for arbitrary $u \in L^2(\mathfrak{X})$ and $\tilde{X} = (x, \xi, t) \in \mathbb{H}_\mathfrak{X}$. When restricted to $\Xi = \mathfrak{X} \times \mathfrak{X}^*$ (which is not a subgroup and should be regarded as a quotient of $\mathbb{H}_\mathfrak{X}$), $\Pi$ becomes a projective representation that will be denoted by $\pi$: it satisfies

$$\pi(X) \pi(Y) = \kappa(X, Y) \pi(X + Y), \quad \forall X, Y \in \Xi.$$
The Wigner distributions defined by \( \pi \) are given by
\[
\mathcal{W}(u,v) := \mathcal{F}(\langle u, \pi(\cdot) v \rangle), \quad u, v \in L.
\]
We used the symplectic Fourier transform
\[
(\mathcal{F} a)(X) := \int_{\Xi} dY e^{-i[X,Y] a(Y)}
\]
and forced it to be \( L^2 \)-unitary and satisfy \( \mathcal{F}^2 = \text{id} \), by a suitable choice of Lebesgue measure \( dY \) on \( \Xi \). Recall that \( \mathcal{W}(u,v) \in \mathcal{S}(\Xi) \) when \( u, v \in \mathcal{S}(\mathcal{X}) \), \( \mathcal{W} : L \times L \to L^2(\Xi) \) is an isometry and extends to a unitary mapping \( \mathcal{W} : L^2(\mathcal{X}) \otimes L^2(\mathcal{X}) \to L^2(\Xi) \). The Weyl pseudodifferential calculus is then a linear isomorphism
\[
\text{Op} : \mathcal{S}'(\Xi) \to \mathcal{B}[\mathcal{S}(\mathcal{X}), \mathcal{S}'(\mathcal{X})], \quad \langle v, \text{Op}(a)u \rangle = \langle \mathcal{W}(v, u), a \rangle. \tag{3.2}
\]
Recall also that \( \text{Op}[\mathcal{W}(u,v)] = \langle \cdot | v \rangle u \) for all \( u, v \in L \), and \( \text{Op} : L^2(\Xi) \to \mathcal{B}_2(L) \) (Hilbert-Schmidt operators) is unitary. For \( a, b \in \mathcal{S}'(\Xi) \), \( a \sharp b \) is the symbol of the operator \( \text{Op}(a) \text{Op}(b) \) whenever this is well-defined and continuous from \( \mathcal{S}(\mathcal{X}) \) to \( \mathcal{S}'(\mathcal{X}) \). Of course, the symbol \( \sharp \) is an extension of the one used in the previous section. The action of \( \text{Op}(a) \) on \( \mathcal{S}(\mathcal{X}) \) or \( L := L^2(\mathcal{X}) \) (under various assumptions on the symbol \( a \) and with various interpretations) is given by
\[
\text{[Op}(a)\text{]}(x) := \int_{\mathcal{X}} dy \int_{\mathcal{X}} d\xi e^{i(x-y)\cdot \xi} a \left( \frac{x+y}{2}, \xi \right) v(y). \tag{3.3}
\]
Consider next the space of operators
\[
\mathcal{B}_u(L) = \{ T \in \mathcal{B}(L) \mid \exists X \to \pi(X)T\pi(-X) \in \mathcal{B}(L) \text{ is norm continuous} \}.
\]
Then \( \mathcal{B}_u(L) \) is a proper \( C^* \)-subalgebra of \( \mathcal{B}(L) \) with the norm given by the operator norm and involution given by Hilbert space adjoint, and it contains the ideal \( K(L) \) of compact operators on \( L \) (see [6 Thm. 1.1]). The representation
\[
\pi \otimes \bar{\pi} : \Xi \to \mathcal{B}_u(L), \quad (\pi \otimes \bar{\pi})(X)T = \pi(X)T\pi(-X)
\]
is then strongly continuous. Let \( \mathcal{B}_u^\infty(L) \) be the space of smooth vectors for this representation. Then \( \mathcal{B}_u^\infty(L) \) is dense in \( \mathcal{B}_u(L) \) [6 Thm.1.1], and consists precisely of those Weyl pseudo-differential operators with symbols in \( BC^\infty(\Xi) \) ([6 Thm.1.2] and [7 Thm. 2.3.7]).

**Lemma 3.1.** The Weyl calculus \( \text{Op} \) realizes an isomorphism between \( \mathcal{H}(\Xi) \) (the R-deformation of \( BC^\infty_u(\Xi) \)) and \( \mathcal{B}_u(L) \). The image through \( \text{Op} \) of \( \mathcal{S}(\Xi) \) is precisely \( K(L) \).

**Proof.** Indeed, recall that when \( a \in BC^\infty(\Xi) \), the norm \( \| a \|_{\mathcal{H}(\Xi)} \) of \( a \) in the Rieffel algebra is given by the norm of the operator \( L_a : \mathcal{S}(\Xi) \to \mathcal{S}(\Xi) \),
\[
L_a(b) = a \# b, \quad \text{where on } \mathcal{S}(\Xi) \text{ one considers the } L^2 \text{-norm (a particular case}
of [15 Prop. 4.15]). Taking \( b = \mathcal{W}(u, v) \), with \( u, v \in \mathcal{S}(\mathcal{F}) \), one gets that
\[
\| L_a \mathcal{W}(u, v) \|_{L^2(\Xi)} = \| v \| \| \text{Op}(a) u \|,
\]
u, v \in \mathcal{S}(\mathcal{F}).
Thus \( \| \text{Op}(a) \|_{\mathcal{B}(\mathcal{L})} \leq \| a \|_{\mathcal{A}(\mathcal{E})} \). On the other hand, denoting by \( \| \cdot \|_{S^2(\mathcal{L})} \) the Hilbert-Schmidt norm, one has
\[
\| L_a(b) \|_{L^2(\Xi)} = \| \text{Op}(a # b) \|_{S^2(\mathcal{L})}
\leq \| \text{Op}(a) \|_{\mathcal{B}(\mathcal{L})} \| \text{Op}(b) \|_{S^2(\mathcal{L})}
= \| \text{Op}(a) \|_{\mathcal{B}(\mathcal{L})} \| b \|_{L^2(\Xi)},
\]
hence \( \| \text{Op}(a) \|_{\mathcal{B}(\mathcal{L})} \geq \| a \|_{\mathcal{A}(\mathcal{E})} \). It follows that the norm of the operator \( L_a \) is in fact equal to the norm of \( \text{Op}(a) \) in \( \mathcal{B}(\mathcal{L}) \). The Rieffel algebra \( \mathcal{B}(\mathcal{E}) \) is the closure of \( \mathcal{B}_\infty(\mathcal{L}) \) in the norm \( a \rightarrow \| a \|_{\mathcal{A}(\mathcal{E})} = \| L_a \|_{\mathcal{B}(\mathcal{L})} \), hence it is isomorphic to \( \mathcal{B}_\infty(\mathcal{L}) = \text{Op}[\mathcal{B}_\infty(\mathcal{E})] \), as stated.

Now the last statement of the Lemma is trivial if we recall that \( \text{Op}[\mathcal{S}(\mathcal{E})] \subset K(\mathcal{L}) \).

## 4 The canonical mappings

**Definition 4.1.** On \( \mathcal{S}(\Xi; \mathcal{A}^\infty) \) we introduce the canonical mappings
\[
[M(F)](X) := \int_\Xi dY e^{-i[X,Y]} \Theta [F(Y)]
\]
and
\[
[M^{-1}(G)](X) := \int_\Xi dY e^{-i[X,Y]} \Theta_{-X} [G(Y)].
\]

To give a precise meaning to these relations, use the (symplectic) partial Fourier transform
\[
\mathfrak{F} \equiv \mathcal{F} \otimes 1 : \mathcal{S}(\Xi; \mathcal{A}^\infty) \rightarrow \mathcal{S}(\Xi; \mathcal{A}^\infty),
\]
\[
(\mathfrak{F}F)(X) := \int_\Xi dY e^{-i[X,Y]} F(Y).
\]

Defining also \( C \) by \([C(F)](X) := \Theta X [F(X)]\), we have \( M = \mathfrak{F} \circ C \) and \( M^{-1} = C^{-1} \circ \mathfrak{F} \).

**Proposition 4.2.** The mapping \( M : (\mathcal{S}(\Xi; \mathcal{A}^\infty), \square, \square) \rightarrow (\mathcal{S}(\Xi; \mathcal{A}^\infty), \circ, \circ) \) is an isomorphism of Fréchet *-algebras and \( M^{-1} \) is its inverse.

**Proof.** The partial Fourier transform is an isomorphism. One also checks that \( C \) is an isomorphism of \( \mathcal{S}(\Xi; \mathcal{A}^\infty) \); this follows from the explicit form of the seminorms on \( \mathcal{S}(\Xi; \mathcal{A}^\infty) \), from the fact that \( \Theta X \) is isometric and from the formula
\[
\partial^\beta \{ \Theta X [F(X)] \} = \sum_{\gamma \leq \beta} C_{\beta \gamma} \Theta X \{ \delta^\gamma \{ \partial^{\beta - \gamma} F(X) \} \}.
\]
With this remarks we conclude that \( M = \mathfrak{F} \circ C \) and \( M^{-1} = C^{-1} \circ \mathfrak{F} \) are reciprocal topological linear isomorphisms.

We still need to show that \( M \) is a \(^*\)-morphism. For the involution:

\[
[M(F)]^*(X) = \left\{ \int_{\Xi} dY \, e^{-i[Y,Y]} \Theta_Y [F(Y)] \right\}^*
= \int_{\Xi} dY \, e^{-i[Y,Y]} \Theta_Y [F(Y)^*]
= \left[ M \left( F^\square \right) \right](X).
\]

For the product: it is enough to show that \( M^{-1} [M(F) \circ M(G)] = F \circ G \) for all \( F, G \in \mathcal{S}(\Xi; \mathfrak{F}^\infty) \). One has (iterated integrals):

\[
(M^{-1} [MF \circ MG])(X) = \int_{\Xi} dY_1 e^{-i[X,Y_1]} \Theta_{-X} \left\{ [MF \circ MG] (Y_1) \right\}
= \int_{\Xi} dY_1 e^{-i[X,Y_1]} \Theta_{-X} \left\{ \int_{\Xi} dY_2 e^{-i[Y_1,Y_2]} \Theta_{(Y_2-Y_1)/2} \left[ (MF)(Y_2) \right] \Theta_{Y_2/2} \left[ (MG)(Y_1 - Y_2) \right] \right\}
= \int_{\Xi} dY_1 \int_{\Xi} dY_2 e^{-i[X,Y_1]} e^{-i[Y_1,Y_2]} \Theta_{-X} \left\{ \Theta_{(Y_2-Y_1)/2} \left[ (MF)(Y_2) \right] \Theta_{Y_2/2} \left[ (MG)(Y_1 - Y_2) \right] \right\}
= \int_{\Xi} dY_1 \int_{\Xi} dY_2 \int_{\Xi} dY_3 e^{-i[X,Y_1]} e^{-i[Y_1,Y_2]} e^{-i[Y_2,Y_3]} e^{-i[Y_1-Y_2,Y_3]} \cdot \Theta_{Y_3/2} \left[ (MF)(Y_3) \right] \Theta_{Y_4} [G(Y_4)]
= \int_{\Xi} dY_1 \int_{\Xi} dY_2 \int_{\Xi} dY_3 \int_{\Xi} dY_4 e^{-i[X,Y_1]} e^{-i[Y_1,Y_2]} e^{-i[Y_2,Y_3]} e^{-i[Y_1-Y_2,Y_3]} \cdot \Theta_{Y_3+Y_2/2} \left[ F(Y_3) \right] \Theta_{Y+Y_2/2} \left[ G(Y_4) \right]
= 2^{4n} \int_{\Xi} dY_1 \int_{\Xi} dZ \int_{\Xi} dY_3 \int_{\Xi} dY_4 e^{-2i[X,Y_1-Y_3]} e^{2i[Y,Z]} e^{-2i[Y_3,Y_4]} \Theta_Y [F(Y)] \Theta_Z [G(Y)] \cdot 
\]

For the last equality we made the substitution \( Y = Y_3 + \frac{i}{2}(Y_2 - Y_1) - X, \ Z = Y_4 + \frac{i}{2}Y_2 - X \). Finally, setting \( Y_3 = X - A, \ Y_4 = X - B \), we get

\[
(M^{-1} [MF \circ MG])(X) = [F \circ G](X) =
= 2^{4n} \int_{\Xi} dY \int_{\Xi} dZ \int_{\Xi} dA \int_{\Xi} dB e^{-2i[A,B]} e^{2i[Y,Z]} \Theta_Y [F(X - A)] \Theta_Z [G(X - B)] .
\]

Remark 4.3. Let us make some comments about how one could modify the definitions above. We are going to need the notation \( [C_\alpha(F)](X) := \Theta_{\alpha X} [F(X)] \), where \( X \in \Xi, \ F \in \mathcal{S}(\Xi; \mathfrak{F}^\infty) \) (or \( F \in L^1(\Xi; A) \)) and \( \alpha \) is a real number. All these operations are isomorphisms and our previous transformation \( C \) coincides
with $C_1$. The traditional composition law in the twisted crossed product is not given by

$$(G_1 \circ G_2)(X) := \int_{\Xi} dY \kappa(X, Y) G_1(Y) \Theta_Y [G_2(X - Y)].$$

The distinction is mainly an ordering matter and it corresponds to the distinction between the Weyl and the Kohn-Nirenberg forms of pseudodifferential theory. Applying $C_1/2$ leads to an isomorphism between the two algebraic structures. So, if we want to land in this second realization, we should replace $M = \mathfrak{S} C_1$ with $M' := C_{1/2} \mathfrak{S} C_1$, leading explicitly to

$$[M'(F)](X) := \int_{\Xi} dY e^{-i[X,Y]} \Theta_{Y+X/2} [F(Y)].$$

5 The $C^*$-algebraic isomorphism

We recall that $\mathcal{H}(\Xi)$, with multiplication $\ast$, has been defined as the R-deformation of the commutative $C^*$-algebra $C_0(\Xi)$ on which $\Xi$ acts by translations. Then $\mathcal{H}(\Xi)$ is an elementary (hence nuclear) $C^*$-subalgebra of $\mathcal{B}(\Xi)$, and $\mathcal{K}(\Xi)$ is dense in $\mathcal{H}(\Xi)$. The Fréchet $^*$-algebra $\mathcal{K}(\Xi; \mathbb{A}^\infty) \equiv \mathcal{H}(\Xi) \otimes \mathbb{A}^\infty$ with the composition law $\Box$ given in (2.8) is dense in the $C^*$-algebra $\mathcal{K}(\Xi) \otimes \mathbb{A}$, that can be viewed (see Remark 2.2 and [16, Prop.2.1]) as the R-deformation of $C_0(\Xi) \otimes \mathbb{A}$ with respect to the action of $\Xi \times \Xi$ composed of translations in the first variable and the initial action $\Theta$ in the second.

This section is mainly dedicated to the proof of the next result:

**Theorem 5.1.** The mapping $M$ extends to a $C^*$-isomorphism $\mathcal{K}(\Xi) \otimes \mathbb{A} \to \mathfrak{S} C_1$.

The following definition (see [15, Def.1.2] and the concept of differential seminorm in [5, Def.3.1]) isolates a situation in which any injective morphism between a dense $^*$-subalgebra of a $C^*$-algebra and another $C^*$-algebra can be extended to a $C^*$-algebraic monomorphism.

**Definition 5.2.** Let $\mathcal{F}$ be a dense Fréchet subalgebra of a $C^*$-algebra $\mathcal{G}$. We say that $\mathcal{F}$ satisfies the Blackadar-Cuntz condition in $\mathcal{G}$ if the topology on $\mathcal{F}$ is given by a family of seminorms $\{p_k\}_{k \geq 0}$ such that $p_0$ is the $C^*$-norm giving the topology on $\mathcal{G}$ and

$$p_k(ab) \leq \sum_{i+j=k} p_i(a)p_j(b), \quad a, b \in \mathcal{F}.$$ 

Tracing back through [5], one realizes that if $\mathcal{F}$ satisfies the Blackadar-Cuntz condition in $\mathcal{G}$, it is a smooth algebra in the sense of [5, Def.6.6]. Actually the more general concept of derived seminorm [5, Def.5.1] involved in the definition of a smooth algebra is meant to model quotients of differential seminorms. Therefore the following result is in fact a particular case of [5, Prop.6.8]:
Proposition 5.3. Assume $\mathcal{F}$ is a dense Fréchet subalgebra of a $C^*$-algebra $\mathcal{C}$ and satisfies the Blackadar-Cuntz condition in $\mathcal{C}$. Then if $\mathcal{D}$ is another $C^*$-algebra and $\Phi: \mathcal{F} \mapsto \mathcal{D}$ is an injective *-morphism, then $\Phi$ is isometric for the $C^*$-norm on $\mathcal{C}$.

We now prove Theorem 5.1.

Proof. The algebra $\mathcal{S}(\Xi) \otimes \mathfrak{A}^\infty$ is a dense subalgebra of $\mathcal{K}(\Xi) \otimes \mathfrak{A}$. As mentioned before, it can be identified to $\mathcal{S}(\Xi; \mathfrak{A}^\infty)$. Proposition 4.2 gives an injective *-morphism $M: \mathcal{S}(\Xi; \mathfrak{A}^\infty) \to \mathcal{A} \rtimes_\theta \Xi$ with dense range. If one proves that $\mathcal{S}(\Xi) \otimes \mathfrak{A}^\infty$ satisfies Blackadar-Cuntz condition in $\mathcal{K}(\Xi) \otimes \mathfrak{A}$, Proposition 5.3 shows that $M$ is isometric for the $C^*$-norm on $\mathcal{K}(\Xi) \otimes \mathfrak{A}$, so it extends to an isomorphism $\mathcal{K}(\Xi) \otimes \mathfrak{A} \to \mathcal{A} \rtimes_\theta \Xi$.

To show the Blackadar-Cuntz condition for $\mathcal{S}(\Xi) \otimes \mathfrak{A}^\infty$, we are going to express it as the space of smooth vectors for a continuous group action in $\mathcal{K}(\Xi) \otimes \mathfrak{A}$.

Using the Schrödinger representation (3.1) of the Heisenberg group $\mathbb{H}_X$ in $\mathcal{L} = L^2(\mathcal{X})$, we consider the strongly continuous representation (by Banach space isomorphisms)

$$
\Delta: \mathbb{H}_X \times \mathbb{H}_X \to \mathbb{B}[\mathcal{K}(\mathcal{L})],
$$

$$
\Delta(\bar{X}, \bar{Y}) T = \Pi(\bar{X}) T \Pi(\bar{Y}), \quad \bar{X}, \bar{Y} \in \mathbb{H}_X.
$$

Notice that $\mathbb{K}(\mathcal{L})$ is an admissible ideal in $\mathbb{B}(\mathcal{L})$, as in [2, Def.3.8]. Also recall that the Weyl-Pedersen calculus for general nilpotent Lie groups $G$, introduced in [11] and developed in [11], particularizes to the usual Weyl calculus if $G = \mathbb{H}_X$ is the Heisenberg group; therefore one can use the results of these articles. It follows from [2, Th.4.6 and Th.3.13] (see also [3] and [11, Th.4.1.4]) that the space $\mathbb{K}(\mathcal{L})^\infty$ of smooth vectors of the representation $\Delta$ is precisely $\text{Op}_{\mathcal{S}(\Xi)}$ and that $\text{Op}: \mathcal{S}(\Xi) \to \mathbb{K}(\mathcal{L})^\infty$ is a topological isomorphism of Fréchet spaces (a restriction of the isomorphism given by Lemma 3.1). Hence $\text{Op} \otimes \text{id}: \mathcal{S}(\Xi) \otimes \mathfrak{A}^\infty \to \mathbb{K}(\mathcal{L})^\infty \otimes \mathfrak{A}^\infty$ is also an isomorphism of Fréchet spaces. Thus, to finish the proof, it will be enough to show that $\mathbb{K}(\mathcal{L})^\infty \otimes \mathfrak{A}^\infty$ satisfies the Blackadar-Cuntz condition in $\mathbb{K}(\mathcal{L}) \otimes \mathfrak{A}$.

We set

$$
\Omega: \mathbb{H}_X \times \mathbb{H}_X \times \Xi \to \mathbb{B}[\mathbb{K}(\mathcal{L}) \otimes \mathfrak{A}],
$$

$$
\Omega(\bar{X}, \bar{Y}, Z) = \Delta(\bar{X}, \bar{Y}) \otimes \Theta Z, \quad \bar{X}, \bar{Y} \in \mathbb{H}_X, \ Z \in \Xi.
$$

It is easy to check that $\Omega$ is a strongly continuous representation and that the space of smooth vectors $\mathbb{K}(\mathcal{L}) \otimes \mathfrak{A}^\infty$ coincides with the (unique) topological tensor product $\mathbb{K}(\mathcal{L})^\infty \otimes \mathfrak{A}^\infty$. The transformations $\Omega(\bar{X}, \bar{Y}, Z)$ are isometric.

It follows that the topology of the tensor product $\mathbb{K}(\mathcal{L})^\infty \otimes \mathfrak{A}^\infty$ is also given by the countable family of seminorms

$$
p_k(\Phi) := \sum_{|\alpha|+|\beta| \leq k} \frac{1}{\alpha! \beta!} \left\| \frac{\partial^{\alpha_1} \partial^{\alpha_2} \partial^{\beta}}{\partial \bar{X}^{\alpha_1} \partial \bar{Y}^{\alpha_2} \partial Z^{\beta}} \left[ \Omega(\bar{X}, \bar{Y}, Z) \Phi \right] \right\|_{\mathbb{K}(\mathcal{L}) \otimes \mathfrak{A}},
$$

\[10\]
where \( \alpha = (\alpha_1, \alpha_2) \in \mathbb{N}^{2n} \) and \( \beta \in \mathbb{N}^{2n} \). When computing on products \( \Phi \circ \Psi \), one has to face the fact that the action \( \Delta \) is not automorphic on \( K(\mathcal{L}) \). Note however that when \( S, T \in K(\mathcal{L}) \), we have

\[
\Delta(\bar{X}, \bar{Y})(ST) = [\Delta(\bar{X}, 0)S] [\Delta(0, \bar{Y})T],
\]

implying for all \( \Phi, \Psi \in K(\mathcal{L}) \otimes \mathfrak{A} \) and all \( (\bar{X}, \bar{Y}, Z) \in \mathbb{H}_\mathcal{L} \times \mathbb{H}_\mathcal{L} \times \Xi \)

\[
\Omega(\bar{X}, \bar{Y}, Z)(\Phi \circ \Psi) = [\Omega(\bar{X}, 0, Z)\Phi] \circ [\Omega(0, \bar{Y}, Z)\Psi].
\]

Then a simple calculation shows that

\[
p_k(\Phi \circ \Psi) \leq \sum_{i+j=k} p_i(\Phi)p_j(\Psi)
\]

for all \( \Phi, \Psi \in K(\mathcal{L})^\infty \otimes \mathfrak{A}^\infty \). One also has \( p_0(\Phi) = \| \Phi \|_{K(\mathcal{L}) \otimes \mathfrak{A}} \), so the proof is finished. \( \square \)

**Remark 5.4.** Let us consider the continuous action \( \beta : \Xi \to \text{Aut}(\mathcal{A} \times_\Theta \Xi) \) given for \( G \in L^1(\Xi; \mathcal{A}) \) by

\[
[\beta_Z(G)](X) := e^{|X, Z|} G(X), \quad X, Z \in \Xi
\]

(this is the dual action in disguise). Then a short computation gives for any \( Z \in \Xi \)

\[
M \circ (T_{-Z} \otimes \Theta_Z) = \beta_Z \circ M,
\]

so actually \( M \) can be seen as an isomorphism of \( C^*-\)dynamical systems. Thus the twisted crossed product \( \mathcal{A} \times_\kappa^{\beta} \Xi \) endowed with the action \( \beta \) can be seen as (an isomorphic copy of) the Rieffel deformation of the \( C^*\)-algebra \( C_0(\Xi; \mathcal{A}) \).

### 6 Applications

One can rephrase Theorem 5.1 by saying that \( \mathcal{A} \times_\kappa^{\beta} \Xi \) is (isomorphic to) the stable algebra of \( \mathcal{A} \). In particular, \( \mathfrak{A} \) and \( \mathcal{A} \times_\kappa^{\beta} \Xi \) are stably isomorphic. Therefore they have identical representation theories (indexed by covariant representations of the system \( (\mathcal{A}, \Theta, \kappa, \Xi) \)), isomorphic ideal lattices and there are canonical homeomorphisms between the corresponding spaces of primitive ideals \([14]\).

We investigate now the interplay between the canonical maps and \( \Xi\)-morphisms. Let \( (\mathcal{A}^j, \Theta^j, \Xi, \kappa) \), \( j = 1, 2 \), be two sets of classical data and let \( \mathcal{R} : \mathcal{A}^1 \to \mathcal{A}^2 \) a \( \Xi\)-morphism, i.e. a \( C^*\)-morphism intertwining the two actions \( \Theta^1, \Theta^2 \). Then \( \mathcal{R} \) acts coherently on \( C^\infty\)-vectors (\( \mathcal{R}[\mathcal{A}^{1, \infty}] \subset \mathcal{A}^{2, \infty} \)) and extends to a morphism \( \mathfrak{R} : \mathfrak{A}^1 \to \mathfrak{A}^2 \) of the R-quantized \( C^*\)-algebras that also intertwines the corresponding actions (see \([12]\)). On the other hand \([12][15]\) another \( C^*\)-morphism \( \mathcal{R}^\times : \mathcal{A}^1 \times_\Theta^{\beta_1} \Xi \to \mathcal{A}^2 \times_\Theta^{\beta_2} \Xi \) is assigned canonically to \( \mathcal{R} \), uniquely defined by

\[
[\mathcal{R}^\times(F)](X) := \mathcal{R}[F(X)], \quad \forall F \in L^1(\Xi; \mathcal{A}^1).
\]
Proposition 6.1. Denoting by $\text{id}$ the identical map on $\mathcal{X}(\Xi)$ and by $M^1$ the canonical map for the data $(\mathcal{A}^1, \Theta^1, \kappa, \Xi)$, one has
\[ R^* \circ M^1 = M^2 \circ (\text{id} \otimes \mathcal{R}). \] (6.1)

Proof. It is enough to compute on $F \in \mathcal{S}(\Xi; \mathcal{A}^\infty)$:
\[
\left[(R^* \circ M^1)(F)\right](X) = R \left[ \int_\Xi dY e^{-i[\xi, Y]} \Theta^1_Y(F(Y)) \right] \\
= \int_\Xi dY e^{-i[\xi, Y]} R \left[ \Theta^1_Y(F(Y)) \right] \\
= \int_\Xi dY e^{-i[\xi, Y]} \Theta^2_Y[R(F(Y))] \\
= M^2[(\text{id} \otimes \mathcal{R})F](X). 
\]

\[\square\]

The next two results have been proved in [17] without asking the skew-symmetric operator $J$ to be non-degenerate (see also [8]). The proofs relying on Theorem 5.1 are very simple:

Corollary 6.2. The $C^*$-algebras $\mathcal{A}$ and $\mathcal{A}^\infty$ have the same $K$-groups.

Proof. Since $\mathcal{A}$ and $\mathcal{A}^\infty \otimes \Xi$ are stably isomorphic, they have the same $K$-theory [4]. On the other hand, by the stabilization trick [12], $\mathcal{A}^\infty \otimes \Xi$ is stably isomorphic to a usual (untwisted) crossed product $(\mathcal{A} \otimes \mathcal{X}) \rtimes \Gamma$ associated to an action $\Gamma$ of $\Xi$ on the tensor product of $\mathcal{A}$ with an elementary algebra $\mathcal{X}$. The vector space $\Xi$ has even dimension, hence by Connes’ Thom isomorphism [4] the $K$-groups of the crossed product coincides with the $K$-groups of $\mathcal{A} \otimes \mathcal{X}$, i.e. with those of $\mathcal{A}$. \[\square\]

Corollary 6.3. The $C^*$-algebras $\mathcal{A}$ and $\mathcal{A}$ are simultaneously nuclear.

Proof. The argument is analogue to the previous one. One must also recall [4, Th.15.8.2] that nuclearity is preserved by stable isomorphism and that the crossed product with a commutative group of a $C^*$-algebra $\mathcal{B}$ is nuclear iff $\mathcal{B}$ is nuclear (for the converse use Takai duality). \[\square\]

If $\mathcal{A}$ is commutative, by Gelfand theory, it is isomorphic (and will be identified) to $C_0(\Sigma)$, the $C^*$-algebra of all complex continuous functions on the locally compact space $\Sigma$ which converge to zero at infinity. The space $\Sigma$ is a homeomorphic copy of the Gelfand spectrum of $\mathcal{A}$ and it is compact iff $\mathcal{A}$ is unital. Then the group $\Theta$ of automorphisms is induced by an action (also called $\Theta$) of $\Xi$ by homeomorphisms of $\Sigma$. We are going to use the convention
\[ [\Theta_X(f)](\sigma) := f(\Theta_X(\sigma)), \quad \forall \sigma \in \Sigma, X \in \Xi, f \in \mathcal{A}, \]
as well as the notation $\Theta_X(\sigma) = \Theta(X, \sigma)$ for the $X$-transform of the point $\sigma$. Let us set $\mathfrak{A} = C(\Sigma)$ for the (non-commutative) Rieffel $C^*$-algebra associated to $C_0(\Sigma)$ by deformation and $C^\infty(\Sigma)$ for the space of smooth vectors under the action $\Theta$.

Remark 6.4. A rather surprising picture follows from a symmetry argument. Using [15, Th.6.5], one gets easily an isomorphism $\mathcal{K}(\Xi) \otimes \mathfrak{A} \cong \mathfrak{A} \rtimes_{\kappa} \Theta$. The complex conjugated cocycle $\kappa$ is defined by the symplectic form $-\{\cdot, \cdot\}$. One usually thinks of $\mathfrak{A}$ as a rather simple $C^*$-algebra, giving after deformation a more complicated one $\mathfrak{A}$. In the commutative case, for instance, we might be surprised that the twisted crossed product $C(\Sigma) \rtimes_{\kappa} \Theta$ decomposes as $C_0(\Xi; \mathcal{K}(\Xi))$.

Remark 6.5. Twisted crossed products with commutative $C^*$-algebras are discussed in [10]. In some situation their primitive ideal space is understood (as a topological space) and this can be transferred by our stable isomorphism to the level of $C(\Sigma)$. By [10, Ex.4.3] for instance, if the action $\Theta$ is free (all the isotropy groups are trivial), then $\text{Prim}[C(\Sigma)]$ is homeomorphic to the quasiorbit space $Q^\Theta(\Xi)$. If in addition $\Theta$ is minimal, $C(\Sigma)$ will be a simple $C^*$-algebra. If $\Theta$ is minimal without being free, the situation is described in [10, Ex.4.11].

Remark 6.6. Assume that $\Xi$ act freely on $\Sigma$. By Theorem 5.1 and [10, Th.4.5], $C(\Sigma)$ is a continuous trace $C^*$-algebra if and only if the action $\Theta$ is proper.

We discuss shortly orthogonality matters. On $\Sigma$ we pick a $\Theta$-invariant measure $d\sigma$. The relationship between the spaces $\mathcal{S}(\Xi; \mathcal{A}^\infty)$ and $L^2(\Xi \times \Sigma)$ depends on the assumptions we impose on $(\Sigma, d\sigma)$. If $d\sigma$ is a finite measure, for instance, one has $\mathcal{S}(\Xi; \mathcal{A}^\infty) \subset L^2(\Xi \times \Sigma)$. Anyhow, the canonical map can be defined independently on $L^2(\Xi \times \Sigma)$.

Proposition 6.7. One has the orthogonality relations valid for $F, G \in L^2(\Sigma \times \Xi)$:

$$\langle M(F), M(G) \rangle_{\Xi \times \Sigma} = \langle F, G \rangle_{\Xi \times \Sigma}. \quad (6.2)$$

Thus the operator $M : L^2(\Xi \times \Sigma) \to L^2(\Xi \times \Sigma)$ is unitary.

Proof. It is enough to note that $M = \mathfrak{F} \circ C$ and to use the fact that $\mathfrak{F}$ and $C$ are isomorphisms of $L^2(\Sigma \times \Xi)$ if $d\sigma$ is $\Theta$-invariant. □

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References

[1] Beltiţă, I. and Beltiţă, D. “Magnetic Pseudo-differential Weyl Calculus on Nilpotent Lie Groups”, Annals of Global Analysis and Geometry 36, no. 3, (2009): 293-322.
[2] Belită, I. and Belită, D. “Smooth Vectors and Weyl-Pedersen Calculus for Representations of Nilpotent Lie Groups” Ann. Univ. Bucharest (math. series) 1 (LIX) no. 1, (2010): 17-46.
[3] Belită, I. and Belită, D. “Continuity of Magnetic Weyl Calculus”, J. Funct. Analysis 260, (2011): 1944-1968.
[4] Blackadar, B. *K-Theory for Operator Algebras*, Springer-Verlag, New York/ Berlin/ Heidelberg, 1986.
[5] Blackadar, B. and Cuntz, J. “Differential Banach Algebra Norms and Smooth Subalgebras of $C^*$-Algebras”, J. Operator Theory 26, no. 2, (1991): 255–282.
[6] Cordes, H.O. “On Pseudodifferential Operators and Smoothness of Special Lie-Group Representations”, Manuscripta Math. 28, no. 1-3, (1979): 51–69.
[7] Folland, G. B. *Harmonic Analysis in Phase Space*, Princeton Univ. Press, Princeton, NJ, 1989.
[8] Kasprzak, P. “Rieffel Deformation via Crossed Products”, J. Funct. Anal. 257, no. 1, (2009): 1288–1332.
[9] Măntoiu, M. “Rieffel’s Pseudodifferential Calculus and Spectral Analysis for Quantum Hamiltonians”, Preprint ArXIV and to appear in Ann. Inst. Fourier.
[10] Packer, J. “Transformation Group $C^*$-algebras: A Selective Survey”, Contemporary Mathematics 167, (1994): 183–217.
[11] Pedersen, N. V. “Matrix Coefficients and a Weyl Correspondence for Nilpotent Lie Groups”, Invent. Math. 118, no. 1, (1994): 1–36.
[12] Packer, J. and Raeburn, I. “Twisted Crossed Products of $C^*$-algebras”, Math. Proc. Camb. Phyl. Soc. 106, (1989): 293–311.
[13] Packer, J. and Raeburn, I. “Twisted Crossed Products of $C^*$-algebras”, II, Math. Ann. 287, (1990): 595–612.
[14] Raeburn, I. and Williams, D. “Morita Equivalence and Continuous-Trace $C^*$-Algebras”, Mathematical Surveys and Monographs, 60, American Mathematical Society (1998).
[15] Rieffel, M. A. “Deformation Quantization for Actions of $\mathbb{R}^{dn}$, Memoirs of the AMS, 506, (1993).
[16] Rieffel, M. A. “Compact Quantum Groups Associated with Toral Subgroups” Representation Theory of Groups and Algebras. Contemp. Math. 145, Providence RI: Am. Math. Soc. 1993: 465-491.
[17] Rieffel, M. A. “K-Groups of C*-Algebras Deformed by Actions of $\mathbb{R}^d$”, *J. Funct. Anal.* 116, (1993): 199–214.

[18] Schweitzer, L. B. “Spectral Invariance of Dense Subalgebras of Operator Algebras”, *Intern. J. Math* 4, no. 2, (1993): 289–317.