Pore constrictions in intervessel pit membranes provide a mechanistic explanation for xylem embolism resistance in angiosperms

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Summary

• Embolism spreading in angiosperm xylem occurs via mesoporous pit membranes between vessels. Here, we investigate how the size of pore constrictions in pit membranes is related to pit membrane thickness and embolism resistance.

• Pit membranes were modelled as multiple layers to investigate how pit membrane thickness and the number of intervessel pits per vessel determine pore constriction sizes, the probability of encountering large pores, and embolism resistance. These estimations were complemented by measurements of pit membrane thickness, embolism resistance, and number of intervessel pits per vessel in stem xylem (n = 31, 31 and 20 species, respectively).

• The modelled constriction sizes in pit membranes decreased with increasing membrane thickness, explaining the measured relationship between pit membrane thickness and embolism resistance. The number of pits per vessel affected constriction size and embolism resistance much less than pit membrane thickness. Moreover, a strong relationship between modelled and measured embolism resistance was observed.

• Pore constrictions provide a mechanistic explanation for why pit membrane thickness determines embolism resistance, which suggests that hydraulic safety can be uncoupled from hydraulic efficiency. Although embolism spreading remains puzzling and encompasses more than pore constriction sizes, angiosperms are unlikely to have leaky pit membranes, which enables tensile transport of water.

Introduction

Xylem sap in vessel-bearing angiosperms crosses numerous intervessel walls from the root to the leaf xylem, depending on the plant size, vessel length, intervessel connectivity, and vessel network topology (Loepfe et al., 2007). It is well known that small openings in the secondary cell wall, which are described as bordered pits, play an important role in hydraulic transport between adjacent vessels, and also in failure of the transport system by gas entry (i.e. embolism) (Choat et al., 2008; Kaack et al., 2019). Since water transport efficiency is tightly related to transpiration and photosynthesis, drought-induced embolism formation can have major implications for plant performance, especially under drought (Li et al., 2016a,b; Sorek et al., 2021). Yet many details about the mechanistic relationship between embolism formation and the anatomical determinants of pits remain unclear.

An angiosperm vessel is estimated to have a median of c. 14 188 intervessel pits, with values for different species varying > 200-fold, from c. 500 pits to > 100 000 (sample size, n = 72 species; Supporting Information Fig. S1, based on data from the literature). Each bordered pit pair has a pit membrane, which is mainly composed of c. 20 nm wide cellulose microfibril aggregates. These pit membranes develop from the primary cell wall and middle lamella, and have a mean diameter of 4.8 ± 2.4 µm (n = 43 species; Jansen et al., 2009, 2011). Before pit membranes become hydraulically functional, hemicellulose and pectin compounds are enzymatically removed (O’Brien, 1970; Herbette et al., 2015; Klepsch et al., 2016). Therefore, fully mature pit

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membranes are nonwoven, fibrous porous media of mainly cellulose, with a thickness between c. 160 nm and 1000 nm (Esau, 1977; Pesacreta et al., 2005; Kaack et al., 2019).

Pit membranes control the immediate entry of gas from neighbouring, embolised conduits, and may become sites of further embolism propagation under persistent drought (Zimmermann, 1983; Brodersen et al., 2013; Brodribb et al., 2016; Choat et al., 2016; Roth-Nebelsick, 2019). Embolism spreading from an embolised vessel to a sap-filled vessel has been described as ‘air-seeding’, but the actual mechanism underlying embolism formation represents one of the major knowledge gaps in our understanding of water transport in plants (Jansen et al., 2018). It is assumed that propagation of drought-induced embolism from one vessel to a neighbouring vessel is affected, among other factors, by pore dimensions of intervessel pit membranes. Here, we use the broader term ‘embolism spreading’ or ‘propagation’ instead of ‘air-seeding’ because embolism spreading includes both mass flow and diffusion of gas across pit membranes. Air-seeding, however, is limited to mass flow only, and embolism formation may not be caused only by mass flow of gas (Guan et al., 2021). Also, ‘embolism’ is used instead of the term ‘cavitation’, because the triggering process leading to embolism is unlikely to be due to the formation of a void by phase transition from liquid to gas, but most likely caused by pre-existing bubbles (Hölttä et al., 2002; Schenk et al., 2017).

Instead of perfectly flat, two-dimensional structures, pit membranes are porous media, with pores that include multiple constrictions, with the respective narrowest constriction in each pore governing the flow of water and gas and, consequently, embolism spreading (Fig. 1; Kaack et al., 2019). Estimates of bottleneck diameters (i.e. constriction sizes) vary from 5 nm to well above 200 nm (Choat et al., 2003; Sano, 2005; Jansen et al., 2009; Hilabrbrand et al., 2016). Part of this variation is caused by sample preparation for imaging by scanning electron microscopy (SEM), which induces a reduction in the pit membrane thickness ($T_{PM}$) of up to 50% during drying, with frequently enlarged pores and cracks (Shane et al., 2000; Jansen et al., 2008; Zhang et al., 2017). Moreover, the challenge is to quantify the size and shape of pit membrane pores using a three-dimensional approach. A three-dimensional model based on data obtained from transmission electron microscopy (TEM) of fresh and shrunken pit membranes indicated a high porosity (i.e. void volume fraction) of 81%, highly interconnected pores with nontortuous, unbending passageways, a lack of dead-end pores, and the occurrence of multiple pore constrictions within a single pore (Zhang et al., 2020). Based on a shrinkage model and gold perfusion experiments, it has been found that constriction sizes in pit membrane pores vary from 5 to $\leq$ 50 nm, with an average diameter of c. 20 nm (Choat et al., 2003, 2004; Zhang et al., 2020). The evidence available suggests that pore sizes are fairly constant for angiosperm species, despite considerable variation in $T_{PM}$. Indeed, pore constriction sizes of c. 20 nm are found both in species with thin (c. 200 nm) pit membranes and those with thick (c. 500 nm) pit membranes (Fig. S2), and there is no evidence for large ($\geq$ 50 nm) pore size differences among species (Zhang et al., 2020). However, Zhang et al. (2020) who conducted experiments using gold particles, recorded only the sizes of the particles that were able to penetrate pit membranes, but did not quantitatively report the penetration depth or the frequency distribution of particle sizes. Could small differences in pore constriction sizes and frequencies in pit membranes explain the relatively variable xylem embolism resistances within angiosperms (Choat et al., 2012)?

Angiosperm species with thick pit membranes were found to be more resistant to drought-induced embolism than species with thin pit membranes (Jansen et al., 2009; Li et al., 2016a,b; Dória et al., 2019; Trueba et al., 2019; Thonglim et al., 2020; Leignonnois et al., 2021). This functional link between $T_{PM}$ and $P_{50}$, which is the xylem water potential corresponding to a 50% loss in maximum hydraulic conductivity ($P_{50}$, MPa), is valid at the interspecific (Li et al., 2016a,b), intragenic (Lens et al., 2011; Plavcová & Hacke, 2012; Scholz et al., 2013) and intraspecific level (Schuldt et al., 2016). Variation in $T_{PM}$ is mainly determined by the number of microfibril layers, with thin pit membranes consisting of fewer microfibril layers than thick pit membranes. Note that the number of layers can be estimated by assuming that cellulose fibres have a diameter of c. 20 nm (Pesacreta et al., 2005), and 20 nm pore spaces between each layer based on gold perfusion experiments (Table 1; Zhang et al., 2020). As such, pit membranes with a thickness between 140 and 1180 nm (Jansen et al., 2009; Li et al., 2016a,b) include between 4 and 30 layers. In our models, bottenlecks in a given pore are formed by the pore constrictions between cellulose fibres within a single layer. Therefore, the number of constrictions within a pore ($N_c$) equals the number of layers (Table 1). Since it is unknown why thin pit membranes are more vulnerable to embolism than thick pit membranes (Jansen et al., 2018), we explore the hypothesis that the likelihood of leaky pores with wide pore constrictions is affected by $T_{PM}$, which could explain why $T_{PM}$ is related to $P_{50}$.

The mismatch between pore size estimations based on colloidal gold perfusion and experimental values of embolism resistance resulted in the hypothesis that a very small percentage of pit membranes might contain large pores (Choat et al., 2003, 2004). These rare pit membrane pores are assumed to cause low embolism resistance. The idea of such leaky, rare pits was further enhanced when variation in $P_{50}$ at an interspecific level was found to decrease with increasing pit membrane surface area in intervessel walls (Wheeler et al., 2005). The ‘pit area hypothesis’ (Sperry et al., 2006), which was later termed ‘rare pit hypothesis’, provided a possible explanation for high vulnerability to embolism, and relied on a largely two-dimensional interpretation of pit membranes (Haque et al., 2007; Christman et al., 2009, 2012; Plavcová et al., 2013). While the rare pit hypothesis follows a plausible mechanism that seems well supported by indirect evidence, it cannot be tested because the existence of a rare pit with a large pore cannot be observed directly, and it is impossible to verify from a statistical point of view. However, a three-dimensional modelling approach to estimating the likelihood of leaky pits is clearly lacking.

The number of layers in a pit membrane may affect the size of the narrowest constriction within a pore that crosses the entire
If embolism propagation is at least to some extent determined by the radius of a pore, the most important dimension of a pore is its minimum diameter, that is, the diameter of the narrowest bottleneck along the pore ($R_{\text{MIN}}$, nm). We can think of this diameter as the ‘effective diameter’ of the pore. The entry of an air–water meniscus or a bubble in a pit membrane is determined by the pore with the largest effective diameter within the pit membrane. Thus, embolism spreading and the minimum hydraulic resistance at the intervessel level are governed by the pore with the largest effective diameter in all pit membranes of a single vessel.

First, we hypothesise that the effective diameter of each pore becomes smaller with increasing $T_{\text{PM}}$ and number of pit membrane layers, as proposed by Kaack et al. (2019) (Hypothesis 1).

This hypothesis is investigated at the individual pit membrane level based on a stochastic pit membrane model. Second, we hypothesise that model-based values of embolism spreading largely agree with embolism resistance measurements for a large number of species (Hypothesis 2). Third, we expect that the probability of having a leaky pit membrane is low at the whole vessel level, and is affected by both $T_{\text{PM}}$ (Li et al., 2016a,b) and the total number of intervessel pits per vessel ($N_{\text{PIT}}$; Hypothesis 3) (Wheeler et al., 2005). The second hypothesis is tested based on experimental data on embolism resistance and anatomical measurements, while two further stochastic pit membrane models are developed to test the third hypothesis. Testing these hypotheses should help us to better understand the functional link between embolism resistance and pit membrane ultrastructure.

### Materials and Methods

#### Pit membrane modelling

To better understand the relationship between $T_{\text{PM}}$ and embolism resistance, we developed three complementary pit membrane models. For reasons of simplicity, we assumed the existence of more-or-less cylindrical pores, which govern...
transport phenomena, and modelled each pore as a three-dimensional object instead of a circular, flat opening (Sperry & Hacke, 2004; Mrad et al., 2018). Following the multi-layered pit membrane model of Zhang et al. (2020), we assumed that each pore penetrates a fixed number of microfibril layers. Each of these layers induces a pore constriction of some random radius (Fig. 1c). An important property of each pore is its effective radius, that is, the radius of the narrowest pore constriction within the entire pore ($R_{MIN}$, nm). We were especially interested in how $R_{MIN}$ was affected by $T_{PM}$ (Hypothesis 1), how modelled embolism resistance based on pore constriction size related to measured embolism resistance (Hypothesis 2), and to what extent the likelihood of leaky pit membranes at the entire vessel level was affected by $T_{PM}$ and/or $N_{PIT}$ (Hypothesis 3).

We developed a first model to estimate pit membrane leakiness at the structural level of a single pit membrane, and two models estimating leakiness at the vessel level. Detailed model descriptions and implementations are provided in Methods S1, S2 and S3.

**Model 1. Pore constrictions in single intervessel pit membranes**

In this model (Fig. 2a), we assumed circular pit membranes with a diameter of 5 µm (estimated from $n = 43$ species, based on Jansen et al., 2009, 2011), each comprising a fixed number of pores, which were defined by a fixed number of pore constrictions. The random radius of each pore constriction was modelled by applying left-truncated normal distributions around mean constriction sizes of 20 nm (Scenario 1) and 100 nm (Scenario 2) in diameter to obtain an upper bound for the number of pores that fit into the membrane, resulting in 12 000 and 1100 pores, respectively.

For $T_{PM}$ values between 140 and 1340 nm, we simulated random diameters of pore constrictions of a whole pit membrane and estimated the probability of encountering at least one pore larger than 35 nm (Scenario 1) or 180 nm (Scenario 2) in a pit membrane, as well as the mean and maximum constriction sizes ($R_{MIN\_mean}$, $R_{MIN\_max}$) of the effective radii $R_{MIN}$. $R_{MIN\_mean}$ and $R_{MIN\_max}$ were compared to the experimental data on embolism resistance with calculated embolism propagation pressures based on a modified Young–Laplace equation.

**Model 2. Leaky pit membranes without hole alignment at the vessel level**

Model 2 investigated the occurrence of leaky pit membranes at the vessel level (Fig. 2b) for $T_{PM}$ values between 50 and 1200 nm and $T_{PIT}$ values up to 400 000. Upper bounds for the probability of encountering at least one large pore, spanning an entire intervessel pit membrane with an effective radius larger than $t$, were calculated based on the probability $P$ of encountering a large hole in any given layer ($P^{N_{PIT}}$) for $P = 0.25$ (Scenario 1) and $P = 0.50$ (Scenario 2). A large pore through the pit membrane was assumed to exist if there was at least one large hole in each layer. At the entire vessel level, the upper bound for the probability of encountering a leaky pit membrane ($P_{LP}$) was given by:

$$P_{LP} = 1 - (1 - P^{N_{PIT}})^{N_{PIT}} \quad \text{Eqn 1}$$

Values of $N_{PIT}$ ranging from 510 to 370 755, with a median value of 14 188, were calculated by dividing the total pit membrane surface area per vessel ($A_{p}$) by the corresponding bordered pit area for 72 species using our original data and data from the literature (Fig. S1; Wheeler et al., 2005; Jansen et al., 2011; Lens et al., 2011; Nardini et al., 2012; Scholz et al., 2013; Klepsch et al., 2016).

**Model 3. Leaky pit membranes with hole alignment at the vessel level**

Model 3 considered the alignment of holes within successive layers by modelling pit membranes as stacks of circular cellulose layers (based on $T_{PIT}$). The pit membranes had a diameter of 5 µm, and each layer was perforated by a fixed number of holes (five in Scenario 1, and 10 in Scenario 2). The holes were randomly

Fig. 2 Three mathematical models to investigate the functional link between pit membrane thickness and effective pore diameters. Model 1 (a) is based on a random number model to estimate the size of the narrowest constrictions of pores that traverse an entire pit membrane. This model is run 10 times following Scenario 1 and Scenario 2, which has small or large pore constrictions for 12 000 or 1100 pores per pit membrane, respectively, and with 4–34 constrictions per pore in 140–1350 nm thick pit membranes. Model 2 (b) examines the probability of large pores in 3000–400 000 intervessel pit membranes within an entire vessel. Pit membranes include up to 34 microfibril layers, assuming either a 0.25 or 0.5 chance of encountering a large hole in a single layer. This model is independent of the actual size of what we consider a large pore, and it does not incorporate alignment of pore constrictions. Model 3 (c) evaluates the probability of encountering pores with a large effective radius at the vessel level (i.e. for 30 000 intervessel pits), with pit membranes consisting of 3–23 microfibril layers, assuming 5 or 10 holes of 200 nm per layer. Alignment of holes was included in Model 3 by simulating random locations of holes in each pit microfibril layer, and requiring minimal overlap between consecutive holes to create a pore. Different shades of grey represent various microfibril layers, and a hypothetical flow path is indicated by the blue lines in (b) and (c).
located in each layer and had radii \( t \geq 100 \) nm. The locations of holes within and across layers were simulated stepwise and repeated \( 10^6 \) times for pit membranes with 3 to 23 layers (corresponding to \( T_{PM} \) values of 100–900 nm; Table 1). A pore only traversed all layers if there existed a sequence of holes that were aligned for each pair of adjacent layers (Fig. 2c). For each scenario, we estimated the probability that at least one hole with \( t \geq 100 \) nm crossed an entire pit membrane and the probability of encountering one large pore in a vessel with 30 000 intervessel pits. Minimal overlap of holes between adjacent layers was assumed to be sufficient for embolism spreading, even if only their edges were overlapping.

**Experimental work**

The three models were complemented by experimental data on embolism resistance (\( n = 31 \) species), \( T_{PM} \) measurements at the centre (\( T_{PM,\text{centre}} \)) and near the edges (\( T_{PM,\text{edge}} \) (\( n = 31 \) species), and the total intervessel pit membrane area per average vessel (\( A_p \), \( n = 20 \) species). The methods applied to obtain these data included well-established, previously published protocols (Wheeler et al., 2005; Sperry et al., 2006; Schulte et al. 2016; Kotowska et al., 2020; Zhang et al., 2020), and they are described in detail in Methods S4. All data included original measurements, except for data retrieved from the literature for embolism resistance for five species, and for \( A_p \) values for four species.

**Statistics and data processing**

Data processing, simulations and statistical analyses were performed using EXCEL, R and MATLAB. Shapiro–Wilk Tests were applied to test for normal distribution. Pearson’s correlation coefficient was used to test for linear correlation. Basic linear and nonlinear regressions were fitted to test whether \( P_{12}, P_{50}, P_{88} \) (Table 2), and the slope of vulnerability curves were related to \( T_{PM} \) or \( A_p \), and could be estimated. For each of the 31 species studied, we estimated embolism resistance by integrating their modelled \( R_{\text{MIN,mean}} \) and \( R_{\text{MIN, max}} \) based on \( T_{PM} \), into the equations for the relationship between \( T_{PM} \) and embolism propagation pressure used in Model 1 (Methods S5). This approach allowed us to compare estimated embolism propagation with experimental values of \( P_{12} \) and \( P_{50} \).

**Results**

How likely are large pores in a pit membrane for a wide range of \( T_{PM} \) values?

Average values of \( R_{\text{MIN}} \) (\( R_{\text{MIN,mean}} \)) are very low in Scenario 1 of Model 1, with values \(< 4.5 \) nm for pit membranes of 150 to 1150 nm in thickness (Fig. 3a). The size of \( R_{\text{MIN}} \) declines considerably with increasing \( T_{PM} \), and the largest values (\( R_{\text{MIN, max}} \)) decrease from radii of \( 20.4 \pm 1.4 \) nm (standard deviation, \( \sigma_{R} \)) to \( 6 \pm 0.6 \) nm (Fig. 3a). \( R_{\text{MIN, max}} \) values are at least 2.4 times and up to 4.9 times larger than the \( R_{\text{MIN, mean}} \) values, decreasing with \( T_{PM} \) (Fig. 3a). The likelihood of having an effective diameter \( \geq 35 \) nm approaches zero (0.00005 ± 0.00009, Fig. 3b) when \( T_{PM} \) is \( > 220 \) nm, or the number of layers \( \geq 6 \), and this therefore only occurs in 0.2 out of 12 000 pores.

For Scenario 2 of Model 1, a similar decline of \( R_{\text{MIN}} \) with increasing \( T_{PM} \) is found (Fig. S3), but with steeper declining likelihood values for large pores with \( T_{PM} \). For a \( T_{PM} \) of 220 nm the likelihood of containing a large pore (defined in Scenario 2 of Model 1 as \( \geq 180 \) nm in diameter) is nearly zero.

How does \( T_{PM} \) relate to measured embolism resistance?

The values of \( T_{PM,\text{mean}} \) vary from 165 nm (± 18 \( \sigma_{R} \)) for *Tilia platyphyllos* to 610 nm (± 79 \( \sigma_{R} \)) for *Olea europaea*, and the median of \( T_{PM} \) is equal to 270 nm (\( n = 31 \) species studied; Table S1). The value of \( T_{PM,\text{centre}} \) is always larger than the value of \( T_{PM,\text{edge}} \), with an average difference of 105 nm, varying from

**Table 2** Overview of the abbreviations used in modelling and experimental parameters used with reference to their units and definitions.

| Modelling acronyms | Units | Definition |
|--------------------|-------|------------|
| \( n \)            |       | Sample size |
| \( N_L \)          |       | Number of microfibril layers in a pit membrane; \( N_L = (T_{PM} + 20)/40 \) |
| \( N_{PIT} \)      |       | Average number of intervessel pits for a vessel with average length and diameter |
| \( P \)            |       | Probability of encountering at least one hole larger than a given threshold in any given layer of a pit membrane |
| \( P_{12}, P_{50}, P_{88} \) | MPa | Xylem water potential corresponding to 12%, 50% and 88% loss of maximum hydraulic conductivity, respectively |
| \( T_{PM,\text{mean}}, T_{PM,\text{centre}}, T_{PM,\text{edge}} \) | nm  | Intervessel pit membrane thickness as measured on transmission electron microscopy (TEM) images of freshly embedded xylem samples; mean value, value around the centre, and value near the edges for a pit membrane (excluding pit membrane annulus) |
| \( \sigma_{R} \)   |       | Standard deviation |

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2.1 nm (T. platyphyllo) to 297 nm (O. europaea), and this difference increases with $T_{PM}$. While there is no relationship between $T_{PM\_centre}$ and the coefficient of variation of $T_{PM\_centre}$ (Pearson’s correlation coefficient, $r$ (29) = 0.09, $P > 0.05$), the range between minimum and maximum values of $T_{PM}$ measured within a tissue strongly increases with increasing $T_{PM}$ ($r$(29) = 0.79, $P < 0.001$). Thus, the species with the thinnest and thickest pit membranes give a $T_{PM}$ range of 80 nm (T. platyphyllo) to 497 nm (O. europaea).

When considering the whole dataset of the 31 species studied, $P_{50}$ values are strongly related to the values of $T_{PM\_centre}$ (Table 3; Fig. 4a), with a logarithmic regression showing an $R^2$ value of 0.57 ($R^2$ (29) = 32.0, $P < 0.001$). An outlier in the $T_{PM}$ vs $P_{50}$ relationship includes Corylus avellana, which shows considerably high $T_{PM}$ values of c. 400 nm for a $P_{50}$ value of −2.02 MPa. Slightly lower correlations are found between the $T_{PM\_centre}$ and

\[ P_{12} (R^2, 29) = 24.4; R^2 = 0.46, P < 0.001, \]
\[ P_{50} (R^2, 29) = 34.2, R^2 = 0.54, P < 0.001; \]
\[ Table 3; Fig. 4a). The $T_{PM\_centre}$ values show a stronger relationship to embolism resistance than $T_{PM\_mean}$ and $T_{PM\_edge}$. The average intervessel pit membrane surface area per vessel ($A_p$, Table S1) shows much weaker correlations with $P_{50}$ and $P_{88}$ than any $T_{PM}$ trait, with the strongest correlation between $A_p$ and $P_{12}$ ($R^2 (2, 18) = 7.75, R^2 = 0.30, P < 0.05$; Table 3).

When limiting our dataset to species with an average intervessel pit membrane surface area per vessel ($A_p$) value, no linear correlation between $P_{12}$, $P_{50}$, or $P_{88}$ and $A_p$ can be found (Fig. 4c), whereas the negative correlations between pit membrane thickness ($T_{PM\_mean}$, $T_{PM\_centre}$, and $T_{PM\_edge}$) and $A_p$ are highly significant (Table 4, Fig. 4b).

Furthermore, we find a power regression with an $R^2$ value of 0.48 between the slope of vulnerability curves and $T_{PM\_mean}$ ($F (2, 29) = 88.4, R^2 = 0.48, P < 0.001$; Table 3), with a decreasing slope being associated with increasing $T_{PM\_mean}$, $T_{PM\_centre}$ and $T_{PM\_edge}$ have a stronger relationship with the slope than $T_{PM\_edge}$ (Table 3). Thus, thicker pit membranes result in more negative $P_{50}$ values and a smaller slope, with $T_{PM}$ mainly affecting $P_{50}$ and less $P_{12}$ and $P_{88}$.

Do modelled and measured embolism resistance correspond to each other for a wide range of $T_{PM}$?

There are clear differences in the estimated pressures that would induce embolism spreading, depending on the surface tension and whether the maximum or mean $R_{MIN}$ values are considered.

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**Table 3** Overview of the $r$- and $R^2$-values between pit anatomical characteristics and embolism resistance.

|          | $P_{12}$ | $P_{50}$ | $P_{88}$ | Slope $T_{PM\_centre}$ | $R_{MIN\_max}$ | $R_{MIN\_mean}$ |
|----------|----------|----------|----------|-------------------------|----------------|----------------|
| $T_{PM\_centre}$ | 0.46*** | 0.57*** | 0.54*** | 0.49*** | 0.79*** | 0.79*** |
| $T_{PM\_mean}$ | 0.44*** | 0.56*** | 0.53*** | 0.48*** | - | - |
| $T_{PM\_edge}$ | 0.31** | 0.41*** | 0.39*** | 0.34*** | - | - |
| $A_p$ | 0.39*** | 0.25*** | 0.22*** | 0.10*** | - | - |

Anatomical measurements include mean values of the intervessel pit membrane thickness ($T_{PM\_mean}$), central pit membrane thickness ($T_{PM\_centre}$), and pit membrane thickness near the annulus ($T_{PM\_edge}$). Embolism resistance is quantified as the xylem water potential values corresponding to 12% ($P_{12}$), 50% ($P_{50}$), and 88% ($P_{88}$) loss of the maximum hydraulic conductivity based on vulnerability curves. The estimation of embolism propagation pressure ($P_e$) is either based on the largest value of $R_{MIN}$ across all pores of a membrane ($P_e R_{MIN\_max}$) or the mean value of $R_{MIN}$ across all pores of a membrane ($P_e R_{MIN\_mean}$), using a modified Young–Laplace equation. $A_p$ total intervessel pit membrane surface area for a vessel with average length and diameter. Only the regressions and correlations that show the strongest relationships are given here. Asterisks indicate significant difference (*, $P < 0.05$; **, $P < 0.01$; ***$P < 0.001$; –, irrelevant and not analysed). For more detailed definitions of the acronyms used in this table, see Table 2.

1 Logarithmic regression.
2 Power regression.
3 Pearson correlation coefficient.
For a surface tension of 72 mN m\(^{-1}\), estimated pressures of embolism spreading, which may largely correspond to \(P_{12}\), are much higher than the measured \(P_{12}\) values, and are even higher than the measured \(P_{50}\) values (Fig. S4). Regression lines of the \(T_{PM} - P_{12}\), \(T_{PM} - P_{50}\) and \(T_{PM} - P_{88}\) relationships, however, fall well within the estimated embolism propagation pressures when a surface tension of 25 mN m\(^{-1}\) (Fig. 4a) is considered. Although absolute values of modelled and measured embolism resistance \((P_{12}, P_{50}, P_{88})\) are not exactly the same, they are significantly related to each other \((r(29) = 0.67, P < 0.01; P_{50} to R_{MIN_mean} and R_{MIN_max}; r(29) = 0.74, P < 0.001; P_{88} to R_{MIN_mean}; r(29) = 0.73, P < 0.001; Table 3; Fig. 5). When \(R_{MIN_max}\) is considered, estimated pressures related to embolism spreading show a small range, with values of 1.2 MPa for a \(T_{PM}\) of 140 nm and up to 2.7 MPa for a \(T_{PM}\) of 758 nm (Fig. 5b), which underestimates embolism resistance \((P_{12}, P_{50}, P_{88})\). Much higher embolism propagation pressures, between 5.6 and 10 MPa, are obtained for estimations based on \(R_{MIN_mean}\), thus overestimating embolism resistance (Fig. 5c,d). There is a clear upper limit to embolism propagation pressure for \(R_{MIN_mean}\) at c. 10 MPa, which is achieved for pit membranes with thicknesses ≥ 600 nm (Fig. 4a).

Modelled embolism propagation pressures based on \(R_{MIN_max}\) are close to measured embolism resistance \((P_{12}, P_{50}, P_{88})\) and are slightly lower than the measured \(P_{12}\) and \(P_{50}\) values (Fig. 5a,c). Estimated embolism propagation pressures based on \(R_{MIN_max}\) are especially close to measured embolism resistance for number of species whose \(P_{12}\) and \(P_{50}\) values are less negative (Fig. 5a,b), while estimated embolism propagation pressures based on \(R_{MIN_mean}\) are much higher than \(P_{12}\) and \(P_{50}\) measurements (Fig. 5c,d).

How likely are leaky intervessel pit membranes at the vessel level?

Based on Model 2, the probability of having a leaky pit membrane in a vessel decreases exponentially with increasing \(T_{PM}\) (Figs 6, S5). For a fixed \(T_{PM}\), the slope of the relationship between \(N_{PTT}\) and the probability of a leaky pore strongly
only 20 species for which we obtained \( T \) values were considered. The extended dataset (\( n = 31 \) species) is provided in Supporting Information Table S1. Since \( T_{PM\text{,mean}} \) is calculated based on the thickness at the centre and the edge, correlations between \( T_{PM\text{,mean}} \) with \( T_{PM\text{,centre}} \) and \( T_{PM\text{,edge}} \) should not be considered and are given in brackets. \( A_p \), total intervesSEL pit membrane surface area per vessel for a vessel with average length and diameter; \( P_{12}, P_{50} \) and \( P_{88} \) xylem water potential values corresponding to a 12%, 50%, 88% loss of maximum hydraulic conductivity, respectively.

| \( A_p \) | \( T_{PM\text{,mean}} \) | \( T_{PM\text{,centre}} \) | \( T_{PM\text{,edge}} \) | \( P_{12} \) | \( P_{50} \) | \( P_{88} \) |
|---|---|---|---|---|---|---|
| 1 | -0.44 | 1 | 1 | 1 | 1 | 1 |
| \( T_{PM\text{,mean}} \) | -0.48* | (1.00**) | 0.94** | 0.58** | 0.92** | 1 |
| \( T_{PM\text{,centre}} \) | -0.32 | (0.96**) | 0.75** | 0.70** | 0.90** | 0.99** |
| \( T_{PM\text{,edge}} \) | -0.33 | 0.72** | 0.85** | 0.83** | 0.83** | 0.83** |
| \( -P_{12} \) | -0.40 | 0.83** | 0.94** | 0.40 0.83** | 0.40 0.83** | 0.40 0.83** |
| \( -P_{50} \) | -0.36 | 0.81** | 0.94** | 0.36 0.81** | 0.36 0.81** | 0.36 0.81** |
| \( -P_{88} \) | -0.36 | 0.81** | 0.94** | 0.36 0.81** | 0.36 0.81** | 0.36 0.81** |

Only 20 species for which we obtained \( A_p \) values were considered. The extended dataset (\( n = 31 \) species) is provided in Supporting Information Table S1. Since \( T_{PM\text{,mean}} \) is calculated based on the thickness at the centre and the edge, correlations between \( T_{PM\text{,mean}} \) with \( T_{PM\text{,centre}} \) and \( T_{PM\text{,edge}} \) should not be considered and are given in brackets. \( A_p \), total intervesSEL pit membrane surface area per vessel for a vessel with average length and diameter; \( P_{12}, P_{50} \) and \( P_{88} \) xylem water potential values corresponding to a 12%, 50%, 88% loss of maximum hydraulic conductivity, respectively. \( T_{PM\text{,mean}} \) mean pit membrane thickness; \( T_{PM\text{,centre}} \) membrane thickness at the centre of the pit; \( T_{PM\text{,edge}} \) membrane thickness at the edge. For more detailed definitions, see Table 2. Asterisks indicate significant difference (*, \( P < 0.05 \), **, \( P < 0.01 \)).

**Discussion**

The results described here indicate that the chance of having large pores in pit membranes decreases strongly with the number of constrictions, and therefore with \( T_{PM} \) (Hypothesis 1). This finding is independent of the actual size of pore constrictions and is supported by a strong relationship between embolism resistance and \( T_{PM} \) (Jansen et al., 2009, 2018; Lens et al., 2011; Scholz et al., 2013; Li et al., 2016a,b; Schultd et al., 2016). Modelled embolism propagation values are significantly correlated with measured embolism resistance values (Hypothesis 2), although they differ in terms of absolute values. There is a good agreement when the dynamic surface tension of xylem sap is taken into account (Yang et al., 2020), but embolism spreading does not seem to represent a function of pore constriction size (\( R_{MIN\text{,max}} \) and \( R_{MIN\text{,mean}} \)) only. Our results also suggest that the likelihood of having a leaky pit membrane within a vessel is extremely low (Hypothesis 3) and is mainly determined by \( T_{PM} \). Overall, pore constrictions provide a mechanistic explanation of why embolism resistance is correlated with \( T_{PM} \), and why pit membranes provide hydraulic safety to angiosperm xylem.

The most narrow pore constriction becomes strongly reduced in size with increasing pit membrane thickness

The three models developed here show a negative correlation between the simulated pore sizes and \( T_{PM} \), which is reflected in a low probability of large pores, both at the level of an individual pit membrane and an entire vessel. Based on Model 1, the chance of having a large pore in a pit membrane thicker than 180 nm is close to zero. Interestingly, the thinnest pit membranes measured in this study (\( c. 165–180 \) nm) are likely to represent a lower limit for \( T_{PM} \), since earlier records of \( T_{PM} \) values below 150 nm (Jansen et al., 2009; Li et al., 2016a,b) are likely artefacts resulting from shrinkage (Zhang et al., 2017, 2020; Kotowska et al., 2020). Thus, angiosperm pit membranes seem to have at least four or five layers of cellulose microfibrils and pore constrictions, which keeps the number of large pores very low for most species. There is a clear conceptual relationship between the thickness of
a fibrous porous medium and the size of the narrowest pore constriction, as also seen for nonwoven, fibrous geotextiles that differ in thickness (Aydilek et al., 2007).

Model 2 suggests that the probability of encountering large pores in intervessel pit membranes of an average vessel follows an exponential pattern over a fairly narrow range of \( T_{PM} \), with critical \( T_{PM} \) values between 200 to 300 nm and 500 to 700 nm for a 0.25 and 0.50 likelihood, respectively, of having at least one hole larger than \( t \) within a single microfibril layer. Although this likelihood cannot be accurately determined due to our limited understanding of embolism spreading and the ultrastructure of pit membranes, we believe that a realistic likelihood would probably lay around 0.25, with 0.50 being too conservative. This assumption is supported by the steeper increase in embolism resistance within the lower \( T_{PM} \) range between 140 and 340 nm than in the higher \( T_{PM} \) range, and by the probabilities of large pores in pit membranes approaching zero for \( T_{PM} > 250 \) nm in Models 1 and 3. We applied a logarithmic regression between \( P_{12} \), \( P_{50} \), \( P_{88} \) and \( T_{PM} \) (Fig. 4a), unlike a linear scaling that was previously suggested (Lens et al., 2011; Li et al., 2016a,b). Interestingly, this logarithmic regression has \( P_{50} \) values approaching 10 MPa for a \( T_{PM} \) of > 1350 nm, which corresponds to the upper physical limit of both xylem water potential and the maximum measured \( T_{PM} \) value for angiosperms (Vilagrosa et al., 2003; Jansen et al., 2009; Kanduč et al., 2020).

A clear limitation of Models 1 and 2 is that the alignment of pore constrictions or holes across all layers of a pit membrane is not considered. Although we do not know whether alignment across different layers is required for mass flow of air across a pit membrane, misalignment could enormously reduce the probability of having a leaky pit membrane and increase tortuosity, because the assumed 20 nm distance between the layers in all three models is low compared to the hole size in Models 2 and 3. Thus, applying stricter criteria to Model 3, such as larger overlap of holes across all layers to obtain a geodesic tortuosity that would be close to 1 (Zhang et al., 2020), would lead to considerably
lower probabilities of leaky pit membranes in a vessel. Nevertheless, even the low amount of overlap applied demonstrates that the chance of having a leaky vessel with 30,000 intervessel pit membranes drastically decreases when $T_{PM}$ values are $c.250$ nm (Fig. 7: 0.25 likelihood scenario of Model 2, Model 3). Since the 0.5 likelihood scenario of Model 2 shows a decrease in leakiness at much higher $T_{PM}$ values than the other models, we consider the predictive value and applicability of this scenario as rather low. It is possible that variation in $T_{PM}$ within a vessel or within the vessel network provides additional chances of leakiness, and small differences in $T_{PM}$ across organs (Kotowska et al., 2020) could influence embolism resistance. Capturing this variation, however, is difficult because measuring $T_{PM}$ may not be straightforward, for instance due to TEM preparation artefacts, aggregation of cellulose fibrils into larger aggregates, and seasonal shrinkage of pit membranes (Schmid & Machado, 1968; Sorek et al., 2021).

The difference between central and marginal $T_{PM}$ questions the modelling assumption of equally spaced cellulose fibres. The slightly negatively charged cellulose fibres may repel each other and are more loosely arranged in the centre than near the edges (Zhang et al., 2016), where the fibres are firmly anchored into the pectin-rich annulus and primary wall. Although the orientation of microfibrils may not be completely random and appears to be directed by a dual guidance mechanism (Chan & Coen, 2020), it seems unlikely that more cellulose fibrils are deposited in the centre than near the annulus, in contrast to torus-bearing angiosperms (Dute, 2015).

How is the size of pore constrictions linked to embolism resistance?

Embolism propagation across pit membranes is strongly dependent on $T_{PM}$ and the number of pit membrane layers, which
control the size of the narrowest pore constriction within a pore. Pit membranes are not different from other nonwoven, fibrous porous media, where the pressure required to force a gas bubble through the medium, the so-called bubble point, is a function of the thickness of the medium and its overall structure (Aydilek et al., 2007). Comparisons of modelled embolism propagation pressures with measurements of \( P_{12}, P_{50} \) and \( P_{88} \) show strong correlations but clear differences in absolute values for most species (Fig. 5), with \( P_{12}, P_{50} \) and \( P_{88} \) values falling between the modelled embolism propagation based on \( R_{\text{MIN,mean}} \) and \( R_{\text{MIN,max}} \) (Figs 4a, 5b,d). As could be expected, \( P_{12} \) values were closest to estimated values based on \( R_{\text{MIN,max}} \). In general, this correlation illustrates that our assumptions in Model 1 are fairly well related to real embolism propagation pressures in plants. Despite the simplicity of the three-dimensional models in this study and the modified Young–Laplace equation, it is remarkable that our modelled \( R_{\text{MIN}} \) values of embolism resistance largely agree with experimental values, without accounting for variables such as dynamic surface tension of xylem lipids, bubble snap-off, changes in temperature, gas solubility, and other xylem anatomical traits. Experimental data on the pressure that is needed to induce embolism in angiosperm xylem shows values between 0.4 and 2 MPa (Choat et al., 2004; Jansen et al., 2009; Christman et al., 2012; Wason et al., 2018), which is more or less in line with \( P_{12} \) values for a wide range of angiosperm species (Bartlett et al., 2016). Moreover, 65% of the species in our study show \( P_{12} \) values that are more negative than \(-2\) MPa, with an average \( P_{12} \) value of \(-2.57\) MPa, which matches the average \( P_{12} \) value of \(-2.65\) MPa of 12 temperate angiosperm species (Schuld et al., 2020).

Embolism propagation across thin pit membranes seems to be determined by pores similar in size to \( R_{\text{MIN,max}} \) due to the high similarity between measurements of \( P_{12}, P_{50}, P_{88} \) with modelled embolism resistance based on \( R_{\text{MIN,max}} \). By contrast, embolism spreading in species with thick pit membranes is affected by pore sizes that can be close to both \( R_{\text{MIN,max}} \) and \( R_{\text{MIN,mean}} \) (Figs 3b, 4a). This finding is in line with the fact that high values of \( T_{\text{PM,mean}} \) show a higher standard deviation than low \( T_{\text{PM,mean}} \) values, while the slope of vulnerability curves becomes lower for species with thicker pit membranes. In addition, the standard error values of \( P_{12}, P_{50}, P_{88} \) tend to increase with increasing embolism resistance (Table S1), that is, higher variation in embolism resistance and lower slopes for embolism resistant species could be linked to increasing variation in \( T_{\text{PM}} \) for species with thick pit membranes. In fact, \( R_{\text{MIN,mean}} \) is expected to provide an upper limit for embolism resistance, since it is unlikely that pore constrictions smaller than average values (i.e. < \( R_{\text{MIN,mean}} \)) will determine embolism spreading. Accordingly, \( R_{\text{MIN,max}} \) offers the least resistance to mass flow of gas moving through a pore space, and provides a good explanation for a lower limit to embolism spreading.

There are three possible reasons why modelled embolism resistance does not match the absolute values of the measured \( P_{12} \) values, and these reasons may not be mutually exclusive. First, the values obtained from Model 1 are based on embolism propagation estimates for a single pit membrane model with a certain thickness, while \( P_{12} \) and \( P_{50} \) values represent hydraulically weighted losses of conductivity at the vessel network level, which is affected by various structural xylem parameters, such as vessel grouping and the ratio of \( T_{\text{PM}} \) to pit membrane area (Levionnois et al., 2021).

Second, estimations based on the Young–Laplace equation should be interpreted with caution due to various poorly known parameters and processes when applying this formula to xylem conduits. Embolism formation in a multiphase environment under negative pressure is highly complicated by, for instance, dynamic surface tension, line tension, the contact angle of the gas–liquid interface within the pit membrane, and highly variable pore sizes (Choat et al., 2004; Law et al., 2017; Schenk et al., 2017; Satarifard et al., 2018; Li et al., 2020; Yang et al., 2020; Zhang et al., 2020). Moreover, pore constrictions and porosity could change if pit membranes become deflected and aspirated against the pit border, which could cause pit membrane shrinkage, reduced porosity and condensivity, or rearrangement of microfibrils (Tixier et al., 2014; Zhang et al., 2017, 2020; Kotowska et al., 2020). However, the mechanical properties of pit membranes remain largely unknown (Tixier et al., 2014).

Third, it is also possible that drought-induced embolism spreading does not occur via mass flow of air–water menisci across intervessel pit membranes, as suggested by the air-seeding hypothesis. The discovery of surfactant-coated nanobubbles in xylem sap provides a complementary mechanism of the mass flow of gas, and highlights the importance of amphiphilic, insoluble lipids associated with pit membranes, and bubble snap-off by pore constrictions (Schenk et al., 2015, 2017, Schenk et al. 2018, Schenk et al., 2021; Kaack et al., 2019; Park et al., 2019). Moreover, diffusion of gas molecules between an embolised and an adjacent vessel could represent an additional way in which gas entry could trigger embolism formation (Guan et al., 2021), which might be largely dependent on \( R_{\text{MIN,mean}} \) and less so on \( R_{\text{MIN,max}} \).

Pit membrane thickness and the number of intervessel pits have different consequences on embolism resistance

We show that \( T_{\text{PM}} \) is a much stronger determinant of the likelihood of leaky pit membranes than the number of pits per vessel (\( N_{\text{PIT}} \)) and the total intervessel pit membrane surface area (\( A_{\text{PIT}} \); Table 4; Fig. 4c). Therefore, our results do not support the rare pit hypothesis (Wheeler et al., 2005; Sperry et al., 2006) but provide a novel view of the relationships between \( N_{\text{PIT}} \) and embolism resistance, and \( A_{\text{PIT}} \) and embolism resistance. Most importantly, our Model 2 shows that \( T_{\text{PM}} \) and \( N_{\text{PIT}} \) affect the likelihood of encountering wide pores differently, with contrasting differences for species with a wide range of \( T_{\text{PM}} \) values. The effect of \( N_{\text{PIT}} \) on vessel leakiness is limited to a narrow range of critical \( T_{\text{PM}} \) values, depending on the assumptions made in Model 2 (Figs 6, S5). The idea that large \( A_{\text{PIT}} \) values lead to a high probability of large pore constrictions in a vessel can be applied to a certain range of \( T_{\text{PM}} \) values. However, the hypothesis that large \( A_{\text{PIT}} \) values increase the probability of large pores in each single vessel is highly doubtful based on the available evidence.
In a general, simplified way, three functional types of intervessel pit membranes can be distinguished based on $T_{PM}$ (Model 2): (1) a thin, risky type, with relatively large pore constrictions, rather low embolism resistance, and little or no reduced embolism resistance for a low number of intervessel pits, (2) a thick and very safe pit membrane type, with narrow pore constrictions, high embolism resistance, and hardly any reduction of embolism resistance for a high number of intervessel pits, and (3) an intermediate pit membrane type, with embolism resistance strongly affected by $N_{PIT}$, where $N_{PIT}$ or other xylem structural traits could potentially be modified during growth to vary embolism resistance in response to the amount of drought experienced. Unfortunately, the exact $T_{PM}$ values that would define these pit membrane types are unclear. Based on leakiness probabilities that are close to zero for $T_{PM} > 250$ nm based on Model 1 and Scenario 1 of Model 3 (Figs 3 and 7), and due to the decreasing slopes of the measured $P_{50}$ values with increasing $T_{PM}$, we roughly estimate that $T_{PM}$ values of the intermediate type are between 150 and 300 nm, which is where the high probability drops from 1 to almost 0 in Fig. 6. This would correspond to 60% of the species in our data set.

Interestingly, embolism resistance in the risky and safe pit membranes (types 1 and 2) is either not affected or is weakly affected by the number of intervessel pits. Since the number of intervessel pits is associated with hydraulic connectivity, which in turn affects hydraulic conductivity and thus efficiency (Loepfe et al., 2007; Mrad et al., 2018), this finding suggests that hydraulic safety can be uncoupled from hydraulic efficiency. In other words, for a given $T_{PM}$, and with considerable variation in $N_{PIT}$, hydraulic conductivity could be affected much more by $N_{PIT}$ than by hydraulic safety (Fig. 6, Fig. S6). Hence, Model 2 provides a novel view on the weak relationship between specific hydraulic conductivity and $P_{50}$ values for many angiosperm species (Hacke et al., 2006; Loepfe et al., 2007; Gleason et al., 2016; Sanchez-Martinez et al., 2020). It would also be interesting to examine whether considerable variation in $T_{PM}$ and $N_{PIT}$ leads to considerable variation in the hydraulic resistance of pit membranes.

The rare pit hypothesis relies on the assumption that for successful embolism propagation, there is always at least one large pore per successive intervessel wall within the xylem network. Our results indicate that the rare pit hypothesis cannot explain embolism propagation at the level of the whole vessel network, since the functional importance of multiple pore constrictions makes it highly unlikely that many vessels contain a large pore for a wide range of $T_{PM}$. In fact, earlier studies that tested the rare pit hypothesis should be considered carefully due to possible artefacts in embolism resistance measurements (Wheeler et al., 2013; Torres-Ruiz et al., 2017). Also, no large pores have ever been found in hydrated pit membranes (Schmid & Machado, 1968; Choat et al., 2003, 2004; Pesacreta et al., 2005; Jansen et al., 2018; Zhang et al., 2020). Finally, primary cell wall development, including the assembly and deposition of cellulose fibrillar aggregates, require highly redundant processes involving the cytoplasm and its cytoskeleton, reducing the likelihood of large gaps in primary cell walls (Chaffey et al., 1997; Oda & Fukuda, 2013; Bourdon et al., 2017; Sugiyama et al., 2017, 2019).

Further progress in understanding embolism spreading in angiosperm xylem will strongly depend on the development of realistic three-dimensional pit membrane and vessel network models (Gaiselmann et al., 2014; Mrad et al., 2018; Li et al., 2020), combined with careful simulations of the chemical and physical interactions within a multiphase environment of gas, water, cellulose, and surfactants.

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Author contributions

LK, MW, LP, HJS, VS and SJ planned and designed the research. EI, ZK, SL, CLT, YZ and BS provided experimental data. LK and MW wrote the manuscript, with input from all co-authors. LK and MW contributed equally to this work.

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Supporting Information

Additional Supporting Information may be found online in the Supporting Information section at the end of the article.

**Fig. S1** Frequency distribution of the number of intervessel pits for a vessel with average dimensions and pit characteristics.

**Fig. S2** Transmission electron microscopy (TEM) images of intervessel pit membranes of different thickness.

**Fig. S3** Results of Model 1, Scenario 2; relationship between $\frac{1}{T_{PM}}$ and the size of a pore constriction.

**Fig. S4** Modelled embolism propagation pressure based on Model 1, with a surface tension of 72 mN m$^{-1}$.

**Fig. S5** Three-dimensional graph based on Scenario 2 of Model 2, with a 0.5 probability of having a large pore in a single pit membrane layer.
Fig. S6 Two-dimensional graph based on Model 2 showing the probability of a large pore in a vessel of up to 400 000 pits per vessel.

Methods S1 Detailed model descriptions.

Methods S2 Model 1, Microsoft Excel function.

Methods S3 R script for Model 3.

Methods S4 Protocols: plant material, xylem embolism resistance, TEM, vessel and pit characteristics.

Methods S5 Equations for the relationship between $T_{PM}$ and embolism propagation pressure used in Model 1.

Table S1 Dataset of the 31 angiosperm species studied, with reference to the anatomical and hydraulic traits measured.

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