The Impact of Microelectronics on High Energy Physics Innovation: The Role of 65nm CMOS Technology on New Generation Particle Detectors

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The High Luminosity Large Hadron Collider (HL-LHC) at CERN will constitute a new frontier for the particle physics after the year 2027. Experiments will undertake a major upgrade in order to stand this challenge: the use of innovative sensors and electronics will have a main role in this. This paper describes the recent developments in 65 nm CMOS technology for readout ASIC chips in future High Energy Physics (HEP) experiments. These allow unprecedented performance in terms of speed, noise, power consumption and granularity of the tracking detectors.

Keywords: CMOS, electronics, micro-electronic design, high energy physics experiments readout, high energy physics, pixel sensor circuit, sub micron, tracking detector

1 INTRODUCTION

Microelectronics has always played a fundamental role in high energy particle physics research. Experiments are fundamentally based on specialized sensors producing an electrical signal upon incident ionizing particle. The signal needs to be amplified, shaped, time-stamped by the front-end electronics and eventually sent off-chip as analog or digital information for further processing (i.e., to boards hosting FPGAs). Information is then analyzed, first locally and then more globally using data from many sensors, or different parts of the apparatus to take decisions on the full event. The core performance of a detector is determined by the chosen sensor, the design and technology of the front-end electronics, together with their interconnection.

HEP experiments have to build more sophisticated detectors in order to exploit the new challenges that experimental particle physics needs to reach. Commercial microelectronics is in constant evolution with technology moving to more and more miniaturized circuits. Benefits are several in terms of logic integration, noise performance, speed, power consumption and low gate capacitance. Although these are major improvements for HEP, great care must be taken because front-end electronics requirements are very specific in scope. Electronics is normally custom-made due to the peculiar requirements of the sensors and of the data acquisition (DAQ) system of the experiment. Within the HEP community the responsibility for the electronics design is taken by a single institute or sometimes collaborations of several institutes are created, both during the R&D phase and the final design. Cost and accessibility are key elements to enable HEP research groups work on new technologies. Last but not least, radiation levels of many HEP experiment are largely above those typical in other fields, including space or military. For this reason, an extensive effort is required to well characterize the radiation resistance of every technology up to the level needed by the HEP experiments.

Figure 1A shows the evolution of industry through the years in developing new technologies and Figure 1B the number of papers published by the HEP community and relevant to R&D work focused on micro-electronics at different nodes. The number of papers does not illustrate the entire
effort invested on R&D in a quantitative way but it does indicates on how and when these efforts evolved through time for the different technology nodes.

Tracking detectors that measure the position of particle are the major drivers for the innovation of electronics in HEP. They are facing the extreme experimental challenges when accelerators are pushed to new physics regime, increasing energy or instant luminosity. Two are the main impacts of this. Firstly, particle density increases accordingly and therefore they need to scale down the dimension of the smallest sensitive area in order to maintain a good separation among particles and a high tracking efficiency. This implies they have to integrate more intelligence in the front-end electronics, while maintaining a low power budget, at least identical to the unit square area with respect to previous detectors in order to preserve the amount of material needed for cooling. Secondly, their electronics is required to survive to the radiation level that increases with integral luminosity. This level reaches its maximum in the inner layers of the experiment that are constituted by tracking detectors. The radiation resistance of MOS has the potential to improve with thinner gate dielectric and therefore might benefit from moving to a further miniaturization, at least for planar technologies.

In the year 2000 the main technology node was 0.25 μm, the so-called quarter micron technology. Its R&D in HEP was pushed primarily by the Large Hadron Collider (LHC) accelerator at CERN, where the first collisions started in 2009. The basic components of the detectors, modules composed of sensors and electronics, started to be built around 2004, and therefore technology choices were to be made at the latest on 2000, leading to the wide usage of quarter micron technology in the LHC. The boost of R&D on 0.25 μm started immediately, as it is visible in Figure 1, while the next innovation in HEP micro-electronics came with some delay with respect to its availability from industry. Only in 2006, a further step down to 130 nm CMOS node in HEP is visible and several detector upgrades of LHC detectors for the phase-1 upgrade were based on this technology.

The next boost has been determined by the HL-LHC program. The new HL-LHC accelerator will allow to reach a new frontier for the particle physics after 2027, as it will determine a factor of five to seven increase in instant luminosity at the LHC and a factor of ten integral luminosity.

The 65 nm CMOS technology node was chosen for detailed exploration for the need of HL-LHC experiments since, with respect to the 130 nm CMOS node, allows for more compact digital logic and memories (×4), higher speed (×2), lower digital power per single circuit (×4) and is a mature but yet still affordable technology. More scaled technologies, would achieve a higher miniaturization of digital logic but a full size ASIC (about $2 \times 2 \text{cm}^2$) required for the HEP experiments, would be far more costly and the necessary R&D would take much more time. Moreover, the voltage gain is shown to decrease in most cases, making an optimal analog design more complex.

Since the attractive solution of a monolithic detector cannot be pursued in the time due for the HL-LHC project, the basic solution for the inner layers of tracking detectors is based on a hybrid pixel detector approach. Also, monolithic detectors are limited in terms of speed for the sensitive part and of radiation hardness, at least in the case of the 130 nm or 110 nm processes that represent the nodes used at present for these detectors.

First studies on 65 nm CMOS technology in HEP began in 2010–2012, about three to four years after it was available from industry. Only in 2013, research has increased substantially in the community. A big boost was the formation of a specific R&D international collaboration, called RD53¹ and of a few projects as [1] or AIDA2020², all bringing together several experts to leverage this technology for HEP. This has allowed the achievement of considerable progress in several areas:

- understanding the radiation hardness of the technology;
- development of basic building blocks (DAC, ADC, Band-Gap, T-sensors, PLL, cable-driver, power-regulators, memory-cells, TDC, high-speed serializer); design of small, low power, fast analog front-ends; development of complex digital circuits; the application of Single Event Upset (SEU) mitigation techniques and their verification on real chips; the integration of large-size ASIC, with complex digital logic, I/O.

¹RD53 Collaboration web site: https://rd53.web.cern.ch/rd53.
²http://aida2020.web.cern.ch/aida2020/, WP4.
This paper describes the basic performance improvements and the new developments of ASIC for HEP that became possible (or will be in near future) thanks to the investment on 65 nm CMOS technology and major R&D work. An overview of applications and state-of-the-art for pixel detectors, tracking, timing detectors, high speed data links and others are presented. Table 1 shows the list of detectors implementing 65 nm CMOS technology for the HL-LHC experiments alongside with the main specifications and requirements.

### 2 RADIATION HARDNESS

In order to use new CMOS technology for HEP front-end electronics, a dedicated study of radiation effects is fundamental since the innermost regions of the experiment can reach unprecedented values of radiation level. In particular, the level of total ionizing dose (TID) in the tracking detectors can reach up to 0.5–1 Grad in the inner layers, values that are two-three orders of magnitude larger than in space applications. Immunity of critical digital circuits to Single Event Effects (SEE), in particular to SEU, are also important for a correct functioning of the electronics.

For the current LHC detectors, the 0.25 μW CMOS technology was characterized with TIDs of a few tens of Mrad, as expected for the innermost layers at the end of the LHC program, before entering to the HL-LHC. This led to the use of special enclosed MOS layout to avoid large leakage currents between source and drain [2], needing special digital cells. The 130 nm CMOS technology was proven to be more radiation hard, so that standard digital cell could be used up to almost 100 Mrad [3] without requiring special layouts techniques.

The first studies on 65 nm technology have shown [4, 5] that is more radiation hard than the previous nodes. Visible degradation begins at 100 Mrad and special precautions have to be taken if going up to 500 Mrad values, with most effects seen for small feature gates, in particular on p-MOS, rather than nMOS. The temperature during irradiation has an effect. More limited degradation is found at low temperature, than at room temperature [5, 6], as shown in Figure 2.

Analogue circuits need to minimize performance variations, therefore they normally make use of large MOS gates to minimize mismatches. In general, small size MOS gates, that are suffering from larger degradation at high total dose, are avoided anyway in analog circuits. Nevertheless, a few design rules, vetoing small feature MOS, have been used to avoid weak

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**Table 1** | Table showing main characteristics of HEP detectors developing micro-electronics in 65 nm. Data taken from paper listed in reference.

| Name         | Detector        | Sensor       | Main feature                  | Purpose                      | Radiation tolerance | Performance                  | I/O                      | Power/Area        | Status                      |
|--------------|-----------------|--------------|-------------------------------|------------------------------|---------------------|-----------------------------|--------------------------|-------------------|-----------------------------|
| MPA          | CMS outer TRK   | Silicon macro pixel | Pixel 100μm x 144μm² | Hit detection, L1 tracking trigger | > 150 Mrad         | ENC=190 e⁻                | 8 x 320 Mbps = 2.56 Gbps L1-trigger, 320 Mbps data | ~220 mW ~90 mW/cm² | Prototype tested; going to production |
| SSA          | CMS outer tracker | Silicon short strip | Strip sensor Cdet 2.5pF/strip | Hit detection | > 100 Mrad          | ENC900e⁻ Peak time = 19 ns | 8 x 320 Mbps = 2.56 Gbps L1-trigger, 320 Mbps data | 6 mW 160 mW/cm² | Prototype under test; going to production |
| RDS3A        | CMS/ATLAS pixel demonstrator | Silicon pixel | Pixel 50 x 50 μm² | Hit detection; bunch identification | > 500 Mrad | ENC = 40–80 e⁻ Threshold <1,200 e⁻ Time walk <22 ns | 4 x 1.28 Gbps | <500 mW/cm² | 13 μW/pix | Produced, tested; choice made for final prototype: One submitted |
| ETROC        | CMS MTD-ETL     | Silicon LGAS | Pixel 1.3 x 1.3 mm² | Time of arrival measurement | 100 Mrad | 40 ps timing | ~160 mW/cm² | 2.7 μW/pix | Small prototype with (4x4) matrix submitted August 2019 |
points. The characterization of circuits before and after irradiation have shown good results, see **Figure 3A**.

For a digital circuit, it is important to know the timing variations, both for combinational and sequential circuits, in order to make a detailed circuit verification that guarantees a good functioning in different corner situations. Measurements have been done on test structures with delay lines and ring oscillators, before and after irradiation, to quantify the effects for different components of the technology digital library. Results are shown in **Figure 3B**.

An extensive work have been done to characterize radiation effects on the technology, in order to introduce reliable models.
FIGURE 4 | Threshold variation for nMOS (A) and pMOS (B) vs gate length: data points and model predictions, from [7].

FIGURE 5 | Schematic representation of the MPA-SSA concept, with two layers silicon sensor identifying an incident particle. Two hits identify a particle, the smaller the curvature of the trajectory, the largest the transverse momentum: this allows to quickly select high-pt L1 by correlating the position of the two hits (A). Schematics of digital logic are shown on (B) From [8].
that allow designers to evaluate the effects on their circuits and to minimize them. Some works are found in Ref. [7], as shown in Figure 4.

### 3 TRACKER DETECTORS

In the outer tracker, the CMS experiment has developed a solution that measures directly in the front-end the charged particle curvature, in order to use it as fundamental information for the L1-Trigger decision. A sandwich of a short silicon strip sensor and a pixel sensor is realized, each with its own front-end ASIC [8]. The pixel sensor is read out by a Macro Pixel ASIC (MPA) that receives the information from the strip sensor ASIC (SSA) simultaneously and has a logic that discriminates interesting particles with high transverse momentum. By using the pixel on one side, a good determination of the 3D direction is possible. Therefore, the MPA can provide particle momentum selection at 40 MHz for the L1-trigger. The schematics of the MPA-SSA is visible in Figure 5.

For the MPA, the main challenges motivating the choice of 65 nm CMOS, are the limitation of the power-budget of the chip of 200–220 mW for a 1.2 × 2.6 cm² ASIC, and radiation hardness. The power is roughly equally distributed among: the analog front-end; the L1 memory; the logic blocks, the clock distribution and the I/O. The digital circuits are powered with 1 V to minimize consumption, while the analog and data I/O drivers use 1.2 V. The block diagram of the MPA is visible in Figure 5. During characterization in laboratory [9], a power consumption of about 184 mW has been measured without radiation. Circuit simulation with expected occupancy at HL-LHC predicts about 5% increase, therefore the MPA fulfills the requirement of 200 mW. The chip has been characterized for total dose and minor degradation have been observed at 150 Mrad.

The SSA [10, 11] needs to maintain low power consumption and to provide high radiation resistance. If it has already been discussed in Section 2 the benefit of using 65 nm CMOS for improved radiation hardness, in term of power it is interesting to compare the SSA with the CBC3 [12], the readout chip used in the silicon strip module at the larger radius of Outer Tracker and made in 130 nm CMOS technology. The power consumption of the SSA analog front-end is about 250 μW/ch while is about 330 μW/ch for the CBC3, therefore the gain here is marginal, and also due to the smaller capacitance of the silicon strip sensor. The real benefit of 65 nm CMOS is shown for the digital circuits. The SSA has eight 320 Mbps output lines, since the strip information have to be sent to MPA without zero-suppression, while the CBC3 provides less information outside the ASIC, with only one 40 Mbps data output. Despite the SSA has 64 times more data output than the CBC3, the digital power is only five times larger. The choice of 65 nm CMOS technology has been fundamental in order to accomplish a module with L1-trigger information obtained with a sandwich of a pixel sensor together with a strip sensor. The total power consumption of the SSA is about 500 μW/ch, for a total of 60 mW, to be compared to about 40 mW of the CBC3.

### 4 VERTEX DETECTORS

A new pixel front-end chip, that guarantees high performance at HL-LHC for the inner tracking layers, must be developed for both the ATLAS and CMS detector upgrades. In this case the use of 65 nm CMOS is pushed to the maximum requirements. The pixel hit rate is up to 3 GHz/cm², the granularity of the pixel must be increased by a factor of five with respect to previous state-of-the-art. The chip must work reliably with an efficiency above 99% when exposed to a harsh radiation environment with a total dose of up to 1 Grad. Meanwhile, low power consumption must also be maintained.

The requirements [13] for the final chip dimension is around 4 cm², with a matrix of nearly 400x400 pixels and more logic in the chip periphery, determining a data readout of up to 3-4 Gbps. Fast output drivers, radiation hard and SEU resistant, must also be developed. Thin planar silicon sensors, with a thickness around 100–150 μm, are foreseen for a large fraction of the pixel detector, and this implies a smaller signal than for Phase-0 and Phase-1 detectors, with a most probable value (MPV) for an incident particle of around 6–7,000 e⁻ after irradiation. This defines high requirements for the analog front-end, aiming very low threshold at a level below 1,200 e⁻ and therefore a very low noise performance for the new chip.

For the pixel innermost layers, the CERN RD53 collaboration has gone through several steps, developing first analog and digital building blocks, all tested for irradiation, then designing and prototyping small scale demonstrators [14, 15] and eventually a large scale prototype, so called RD53A [16], that is 2 cm² in size. Table 2 shows the evolution of these prototypes, with their different features. A high granularity, with a pixel dimension of 50 μm x 50 μm was successfully achieved, determining a big challenge for both analog and digital electronics of the matrix. Each shares 50% of area and power.

Three analog front-end designs with different architectures have been realized [17]. All were proven to satisfy the requirements of RD53A, working at thresholds of about 1,200 e⁻ with a noise occupancy at 10⁻⁶ level, a time-walk smaller than 25 ns and a charge deposition measurement with Time-of-Threshold (ToT) technique. Some results are shown in Figure 6. For ToT digitization four bit counting has been found to guarantee good spatial resolution and therefore has been preferred to higher values (5 or 6-bits) in order to determine small pile-up and to limit data readout.

The digital matrix logic, has to handle the prospected high data rates for long L1-trigger latency (12.5 μs, more than double the value at LHC). It has been chosen to store locally the hit information within a region shared by several pixels. This architecture sends only the selected data to chip periphery, removing locally more than 98% of stored data. This results in a faster and low power data transfer to the chip periphery. Storing information from multiple hits of the same cluster translates to significant savings in storage resources; sharing of latency buffers by grouping pixels into so-called regions leads both to compact
### TABLE 2
Table with the different characteristics of FE65P2, CHIPX-FE0, RD53A demonstrators, produced on Sept-2015, July-2016 and Aug-2017 respectively From [16].

| Characteristics \ Demonstr. | FE65P2 | CHIPX65-FE0 | RD53A |
|-------------------------------|--------|-------------|-------|
| Pixel matrix                  | 64 × 64| 64 × 64      | 400 × 190 |
| Matrix organization           | (2 × 2) analog islands; pixel regions; (4 × 64) COREs | (2 × 2) analog islands; pixel regions; (4 × 64) COREs | (2 × 2) analog islands; pixels regions; (8 × 8) COREs |
| Pixel regions                 | (2 × 2) pixels; distributed data buffer; trigger matching | (4 × 4) pixels; centralized data buffer; trigger matching | (2 × 2) or (4 × 4) pixels; trigger matching |
| VFE                           | VFE-1  | VFE-2        | VFE-1, VFE-2 |
| Analog-digital isolation      | Analog triple well; digital triple well | Analog triple well | Analog triple well; Digital triple well |
| Signal digitization           | 4-Bits | Binary or 5-bits | 4 or 8 bits |
| Building blocks               | Few not RD63 | BandGap, DAC, ADC SER, sLVS-Tx/Tx | BandGap, DAC, ADC ana-buffer, PON-reset, Shunt-LDO, sLVS-Tx/Rx, cable driver, PLL-CDR, temp sensor |
| Bias-distribution             | DAC, current mirror for all pixels (single stage mirroring) | DAC, current mirror for column bias-cells (double stage mirroring) | DAC, current mirror for double-column bias-cells (double stage mirroring) |
| Radiation hard design         | Analog  | Analog       | Analog and digital |
| Powering                      | Standard | Standard     | Serial-powering |

**FIGURE 6** ENC vs current consumption for three different RD53A analog front-ends (A,B,C) at different irradiation doses From [26].
circuitry and low power. The exact way in which hit data are stored must be optimized for new requirements, taking into account the cluster size distribution, which in turn depends on the sensor type and the location in the detector, and on the physics input. Different solutions in RD53A have been evaluated with regions of $2 \times 2$ or $4 \times 4$ pixels. Simulation studies of hit losses determined by the local storage in the pixel region for the two different architectures are shown in Figure 7, from [18, 19].

Measurements [20] of readout RD53A efficiency for different trigger latency vs single hit flux have been done. An X-ray source (60 kV Chrome target with a kα peak of 5.4 keV) generates a single hit physical signal to the silicon pixel sensor bump-bonded to RD53A. A calibration signal is injected to the pixel and, if not in pile-up with another hit, is eventually stored into a latency memory position that could eventually be overwritten by the continuous X-ray flux. The efficiency of getting the calibration signal measures the readout efficiency due to the limited memory storage and decreases with an increase of hit flux or trigger latency. Measurements results are shown in Figure 8: at 3 GHz/cm² there is an overall efficiency of 99% and 97.5% for 3 and 12.5 µsec latency. Simulation predictions are also shown in Figure 8 and well agree with data for X-ray flux, validating the model used. In Figure 8 is also displayed the simulation results for charged particle, that presents clustered hits: a prediction of 1% inefficiency is obtained at 3 GHz/cm² and 12.5 µsec latency, satisfying the requirements. The main cause of inefficiency is the four bit digitization of the very-front end signal, based on the Time-Over-Threshold technique, based on a 40 MHz clock for RD53A: with an average of roughly 8 counts, it implies a constraint in the analog dead-time of 200 ns that at the max flux rate of 3 GHz/cm² implies 1% inefficiency. For the final design an equivalent 80 MHz digitization have been chosen to reduce this inefficiency by half.

The schematics of the RD53A periphery is shown in Figure 9 for the digital architecture and the connection to the I/O, sending data with four lines of 1.28 Gbit/s (pixel data and slow-control data) and receiving triggers and synchronous commands in a single differential line with 160 MHz clock. The need of a large bandwidth for data output, minimizing the amount of lines of readout, demands high frequency operation that can only be met by the 65 nm technology at these high radiation levels.

In terms of radiation hardness, RD53A was found to feature minor degradation at levels up to 500 Mrad for the digital circuitry and up to 800 Mrad for the analog.

The baseline prototype for CMS and ATLAS (called RD53B) is the evolution of RD53A, with improvements in many items, on the analog front-ends, digital architecture, I/O and the inclusion of several other features requested by the experiments, like data compression. A first version of RD53B has been submitted on March 2020 for the ATLAS experiment and is currently under test. A new version of the CMS pixel readout chip, again within the

![Figure 7](image1.png)

**FIGURE 7** | Hit losses determined by the dimension of the local pixel region data storage are shown for simulation studies in the case of $2 \times 2$ (A) and $4 \times 4$ (B) architectures studied From [18, 19].

![Figure 8](image2.png)

**FIGURE 8** | Efficiency vs hit rate flux generated by X-ray, that generate single hit clusters: data well fit the simulation prediction. Simulation prediction for clustered hits coming from particles are also shown, see text From [20].
platform of RD53, is going to be submitted for production in the first quarter of 2021.

5 TIMING LAYERS

The ETROC [21] is the readout ASIC in 65 nm technology foreseen as front-end for the Low Gain Avalanche Detector (LGAD) sensor pixel of the MIP Timing Detector (MTD) detector of CMS. The main design challenge is to extract precision timing information of 50 ps from the sensor with a capacitance of about 3.4 pF while keeping the power consumption low. The LGAD itself determines about 30 ps, while the system clock is designed to be precise at 15 ps, thus ETROC should perform at 40 ps.

The basic architecture is shown in Figure 10.
The small prototypes of the ETROC have been realized, with a single channel containing the preamplifier and discriminator (ETROC0) and a 4 × 4 matrix including also the TDC (ETROC1). Data from a test beam with a ETROC0 chip connected to an LGAD sensor show a timing precision of 30–40 ps that is very encouraging. An ETROC1 chip bump bonded to an LGAD sensor has started testing in Fall 2020. Estimation from these prototypes of the power consumption per pixel of ETROC is 2.7 mW maximum, demonstrating a potential lower power with respect to 130 nm. It is anyway still too early to understand the full gain of using 65 nm (CMS design) with respect to 130 nm (ATLAS design). The request of radiation hardness is of 100 Mrad, that for 65 nm CMOS should not be a problem, but needs confirmation for this specific application.

6 HIGH SPEED DATA LINKS

The upgraded detectors for HL-LHC have higher granularity, up to a factor of five or more with respect to present, while the L1 trigger rate foreseen from experiments will increase of an order of magnitude, from 100 kHz to almost 1 MHz. This implies an unprecedented data bandwidth to be read out by the front-end electronics. This has motivated the development of high speed links close to the detector for the necessary I/O, that should eventually guarantee sufficient radiation hardness and low power.

The low power GigaBit Transceiver (lpGBT) is a radiation tolerant ASIC that implements multipurpose high speed bidirectional links, with 2.56 Gb/s downlink and 5.12 or 10.24 Gb/s as uplink, providing Trigger, Data and Slow control information. It can receive several electrical up-links (E-links) from front-ends from 160 to 1,280 Mbps and send information with down-link with 80 to 320 Mbps. A schematics of the lpGBT is visible in Figure 11. A good reference can be found in [22] on the architecture and performance of the core part of this ASIC, as the PLL for data output and the CDR for the data input. The lpGBT is

TABLE 3 | Evolution of GBT to lpGBT, a comparison table of the essential specifications.

| Specification \ ASIC | GBTx | IpGBTx |
|----------------------|------|--------|
| Technology           | 130 nm | 65 nm |
| Up-links             | 4.8 Gbps | 5.12 or 10.24 Gbps |
| Down-links           | 4.8 Gbps | 2.56 Gbps |
| Up E-links from front-end [Mbps] | 80/160/320 | 160/320/640/1280 Mbps |
| Down E-link to front-end | 80/160/320 | 80/160/320 Mbps |
| Radiation hardness   | > 100 Mrad | < 200 Mrad |
| SEU protection       | TMR, FEC | TMR, FEC |
| Phase resolution     | 50 ps | 50 ps |
| Jitter               | < 15 ps@40 MHz | < 5 ps @ 40 MHz |
| Power                | 2.2 W | 320–490 mW (depending on #/type links) |
| Area                 | 12×12 mm² | 9×9 mm² |

[22]https://ep-dep.web.cern.ch/rd-experimental-technologies.
designed in 65 nm CMOS technology and represents the evolution of the GBT made in 130 nm. A comparison on the essential technical specifications of lpGBT and GBT is shown in Table 3 showing the considerable improvements made in high speed data links for HEP experiments with the use of 65 nm CMOS technology.

7 CONCLUSION

Progress in HEP experiments depends on innovation of detectors and the scaling of micro-electronics plays a fundamental role on that. To confirm this, this paper has shown how the use of 65 nm CMOS technology as a new tool for HEP has granted in many ways to reach new goals for the HL-LHC program. It allows integration of more local intelligence, to increase the detector granularity, or to achieve power reduction. It can improve the signals quality from the sensitive elements, as very low noise or fast timing. It can move larger amount of data or faster commands using high frequency clocks.

Future physics programs like ILC, CLIC, FCC will see the 65 nm CMOS as a baseline micro-electronics technology in several more detectors. Reduction in prices over time is normally expected for technologies, and this will make 65 nm CMOS even more attractive and accessible.

Other technology nodes are going to take the today role of 65 nm for HEP in the incoming years, and 28 nm CMOS is a good candidate. This node, as shown in Figure 1, is already being used in a few HEP R&D projects for small size ASIC. Moreover, radiation characterization [23] is quite advanced. CERN has recently launched a new R&D on experimental technologies, and 28 nm and FinFET CMOS technologies are among the key elements. This will make possible to establish in more solid grounds the access to the HEP for new innovative detectors.

AUTHOR CONTRIBUTIONS

The author confirms being the sole contributor of this work and has approved it for publication.
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