Evaporated MoO$_x$ as General Back-Side Hole Collector for Solar Cells

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Abstract: Substoichiometric molybdenum oxide (MoO$_x$) has good potential as a hole-collecting layer in solar cells. In this paper, we report on the application of ultrathin evaporated MoO$_x$ as a hole collector at the back side of two distinct photovoltaic technologies: polymeric and silicon heterojunction (SHJ). In the case of polymer solar cells, we test MoO$_x$ as a hole transport layer in devices with inverted architecture. The higher transparency of the MoO$_x$ film, compared to the commonly used poly(3,4-ethylenedioxythiophene):poly(styrenesulfonate) (PEDOT:PSS), allows an enhanced back reflected light into the photoactive layer, thus boosting the photogeneration, as found from the illuminated $J$-$V$ and external quantum efficiency (EQE) curves. The higher fill factor (FF) of the MoO$_x$-based device also suggests an improved charge collection efficiency compared to the cells with PEDOT:PSS. As for SHJ solar cells, we show that MoO$_x$ offers the means for dopant-free hole collection with both p-type and n-type Si wafers. In the present comparison over planar test structures with Ag back reflecting electrodes, we observe an efficiency gain of approximately 1% absolute against a baseline with a conventional p-type amorphous silicon hole collector. The gain is linked to the increased $V_{OC}$, which is likely due to the reduced recombination at the Si wafer.

Keywords: carrier-selective contacts; molybdenum oxide; polymer solar cells; inverted architecture; silicon heterojunction solar cells

1. Introduction

Transition metal oxides, and substoichiometric molybdenum oxide (MoO$_x$, $x < 3$) in particular among them, have attracted considerable interest for application as a supportive layer for charge collection into solar cells of various technologies, from silicon-based to polymeric to the novel perovskite solar cells [1–12]. Solar cells fundamentally consist of a semiconductor light absorber equipped with selective contacts for charge-carrier separation and extraction, complemented by proper electrodes for external connection. Based on its high work function (up to 6.9 eV) and large bandgap (approximately 3 eV) [13], MoO$_x$ is an excellent candidate as a hole collector with potentially low optical losses independent of the specific photovoltaic technology.

Thermal evaporation is often used for the deposition of thin MoO$_x$ films, making it a viable and cost-effective choice. A drawback is the sensitivity of the material to air and moisture [14], which means that exposure to the environment should be avoided. The issue is emphasized with evaporated films due to their oxygen deficiency. However, oxygen vacancies are to some extent desired, since they enable the charge carrier transport [1].
The material has been more commonly proposed as a front side (illumination side) hole collector, to fully take advantage of its transparency. This type of application, although promising, may entail some technical difficulties for some device architectures. Often, a capping transparent conductive electrode, typically deposited by sputtering, is indeed needed, and it has been shown that Ar plasma exposure and transparent conducting oxide (TCO) overlayer deposition induce coloration of the material and then parasitic light absorption [3,15]. While countermeasures are under investigation [15], we here explore the application of the material as a rear side hole collector, where this issue can be avoided. In this case, the use of a TCO is not mandatory, and a full area metal electrode can be deposited over the MoO$_x$ layer. To demonstrate the wide applicability of the approach, we have here considered two distinct photovoltaic technologies: polymeric and silicon heterojunction (SHJ), where both the fabrication processes and the mechanisms behind the device operation are very different. We report on the application of evaporated MoO$_x$ layers to these types of solar cells and on the associated benefit.

1.1. Focus on Polymer Solar Cells

Polymer solar cells (PSCs) have attracted intensive investigation in the last three decades due to their potential advantages such as mechanical flexibility, light weight, large-area fabrication, and low manufacturing cost by using roll-to-roll printable techniques [16–20]. In order to compete with inorganic devices in the commercialization of real-world applications, high power conversion efficiency (PCE) and prolonged lifetime are of crucial importance. Recently, PCE higher than 18% [21] has been achieved for single-junction devices, but long stabilities are still an issue, although important progress has been recently achieved [22].

The typical PSC design (standard structure) consists of a transparent front electrode covered by a thin hole conducting layer, a photoactive layer (typically a blend of a semiconducting donor polymer and a fullerene derivative acceptor forming a bulk-heterojunction (BHJ) [16]), an electron transport layer, and a rear metal contact. This structure suffers from device stability due to the deterioration of the metal contact (Ca/Al) and hole transport layer (HTL), which is usually poly(3,4-ethylenedioxythiophene):poly(styrenesulfonate) (PEDOT:PSS). The development of the so-called inverted device structure (i.e., a PSC where the polarity of the electrodes is reversed compared to the standard one) has contributed to significantly improving the lifetime of the devices [23,24].

In both the architectures, poly(3,4-ethylenedioxythiophene):poly(styrenesulfonate) (PEDOT:PSS) is the most common material used as an HTL. This material raises some concerns when applied into inverted PSCs. In particular, the processability of PEDOT:PSS (dispersion in water and post deposition annealing) limits its application due to deterioration of the underlying blend, since a lot of blend materials, used in devices with high PCE, suffer from humidity or annealing treatments. Thanks to their favorable optoelectronic properties, metal oxides (such as MoO$_x$) have emerged as possible substitutes of PEDOT:PSS. Solution-processed MoO$_x$ has been successfully used as an HTL in both normal and inverted PSCs [25–27], although application in the inverted device is still a challenge due to processing issues (i.e., deterioration of the blend film). These issues could be avoided when applying thermally evaporated MoO$_x$, which is a promising material as HTL in inverted PSC due to the favorable energy alignment between its work function and the HOMO level of the commonly used blends. For these reasons, we have investigated the use of thermally evaporated MoO$_x$ as an HTL in PSCs with inverted structure and compared their performance with reference devices employing PEDOT:PSS.

1.2. Focus on Silicon Heterojunction Solar Cells

SHJ solar cells are among the most efficient cells, with a record efficiency of 26.7% reached in the interdigitated back contact configuration [28]. Conversion efficiencies above 25% have been also demonstrated in the more conventional architecture featuring contacts on both the front and back sides [29], which is the design here considered. The device is based on a crystalline silicon (c-Si) wafer as a light absorber, which is passivated by ultrathin intrinsic hydrogenated amorphous silicon (a-Si:H) providing the chemical passivation of Si dangling bonds at the c-Si surface. Thin suitably
doped silicon-based layers, mostly a-Si:H, nc-Si:H, or nc-SiO\textsubscript{x}:H, are applied as selective contacts (p-type doped films for hole collection and n-type for electron collection), where doping is used for shifting the Fermi level close to the valence or conduction band of the absorber, thus achieving hole or electron selectivity, respectively. In the top/rear contact design, the two films are separately deposited over the two passivated surfaces. Then, a TCO layer is applied as an antireflective front electrode, complemented by a metal grid, while the back side is covered by a full area metal-based electrode.

The doped layers are fabricated by plasma-enhanced chemical vapor deposition (PECVD) with a gas precursor mixture that includes toxic dopant gases. Doping of the thin film silicon layers requires delicate deposition optimizations and, especially with p-type doping, induces the formation of recombination-active defects. Thus, the possible substitution of these layers with dopant-free counterparts, in addition to fabrication by a simple low-cost technique such as thermal evaporation, is very attractive. In recent years, evaporated MoO\textsubscript{x} has been indeed proposed as an alternative hole collector in SHJ, with this layer applied at the front side of devices based on n-type c-Si wafers, obtaining conversion efficiencies up to 23.5% [7]. This is the most common configuration for this wafer type, which is also called front emitter configuration (minority carrier collector placed at the illuminated front side), where in principle, MoO\textsubscript{x} is advantageous also from the optical point of view if suitable measures are taken to avoid the mentioned deterioration of the material when applying the capping TCO.

As a complementary approach for SHJs with a dopant-free hole collector, we have explored the use of MoO\textsubscript{x} over the rear side. We propose the application of the layer to both n-type and p-type wafers, which are both interesting because the former have so far demonstrated the highest solar cell efficiencies and the latter are still the main choice of the PV manufacturing industry. With the commercially relevant p-type wafer, the implementation of MoO\textsubscript{x} at the rear side corresponds to the standard cell configuration, since the layer functions as a collector of majority carriers, while the n-type Si-based emitter layer is at the front side. Particularly interesting is the application of the MoO\textsubscript{x} layer at the rear of n-type wafers, as it makes also for an unconventional cell design, since MoO\textsubscript{x} functions as an emitter (minority carrier collector). This second configuration (rear emitter cell design) is in itself a promising concept that has recently been under thorough investigation, as it provides relaxed constraints on the optoelectronic properties of the front TCO [30–32].

2. Materials and Methods

2.1. Single-Layer Fabrication and Characterization

MoO\textsubscript{x} layers were deposited onto Eagle XG glass substrates for optical characterization. The amorphous MoO\textsubscript{x} layers were grown by the thermal evaporation of high-purity stoichiometric powder (99.95%) at a rate of 1 Å/s out of a tungsten boat at a base pressure below 4 × 10\textsuperscript{-7} mbar. During the evaporation process, no intentional heating was applied to the substrate, which remained close to ambient temperature. The thickness and optical constants of the samples were evaluated using a variable angle ellipsometer (VASE, J.A. Wollam Co., Lincoln, NE, USA) by acquiring data over three angles of incidence and complementing it with transmittance spectra measured on the same spot. The work function (WF) of the material evaporated on an n-type cSi substrate was measured using a SKP5050 Scanning Kelvin Probe system by KP Technology (Wick, UK) [33,34].

2.2. Device Fabrication

For the fabrication of the PSCs, we used the following structure: glass/ITO/ZnO/photoactive layer/HTL/Ag, where the photoactive layer is a BHJ of poly[4,8-bis(5-(2-ethylhexyl)thiophen-2-yl)benzo[1,2-b:4,5-b’]dithiophene-2,6-diyl-alt-(4-(2-ethylhexyl)-3-fluorothieno[3,4-b]thiophene)-2-carboxylate-2,6-diyl]] (PTB7-Th) and [6,6]-phenyl C\textsubscript{71} butyric acid methyl ester ([70]PCBM), and the HTL is MoO\textsubscript{x} or PEDOT:PSS. The ZnO precursor was prepared by the sol–gel process as described elsewhere [35]. The ZnO films (40 nm) were obtained by spin coating the solution at 4000 rpm for 60 s on top of
patterned and cleaned glass/ITO substrates (Delta Technologies, sheet resistance of 10 Ω/square). The films were annealed at 150 °C for 5 min. The active layer (90 nm) was realized by dissolving PTB7-Th (Solarmer Inc., Irwindale, CA, USA, 15 mg/mL) and [70]PCBM (Solenne BV, Groningen, The Netherlands, 22.5 mg/mL) in a mixed solvent of 1,2-dichlorobenzene:1,8-diiodoctane (97%-3% by volume) and spin coated on top of the ITO/ZnO structure. A 5 nm MoO₃ layer and a 100 nm Ag layer were sequentially thermally evaporated, without breaking the vacuum, through a shadow mask to form a top anode. For the cells with PEDOT:PSS (Clevios AL 4083, Heraeus, Hanau, Germany), a 40 nm-thick layer was deposited as described elsewhere [36] before Ag deposition. The active area of the device was 22 mm².

For the fabrication of Si heterojunction solar cells, we used flat float-zone c-Si <100> wafers, both p-type and n-type, with resistivity of 1–5 Ω·cm (Topsil GlobalWafers A/S, Frederikssund, Denmark). The fabrication sequence started with RCA cleaning of the Si wafers followed by native oxide removal through dipping into 2% HF solution. Then, the two wafer sides were passivated with 5 nm-thick intrinsic a-Si:H. Afterwards, the front side was covered with the electron collector, which is made of highly transparent n-type nanocrystalline silicon oxide (nc-SiOₓ:H) with a thickness of 20 nm, and the antireflective transparent front electrode, which is made of 80-nm sputtered ZnO:Al completed with an Al collection grid. The area of the solar cells was defined as 2 x 2 cm² at the stage of the front electrode fabrication by using a shadow mask. The hole collector was applied to the back side. Two materials were considered in this case: 20 nm-thick conventional p-type a-Si:H and MoO₃ (10 nm-thick, unless otherwise stated). Finally, the back electrode was applied by evaporating a full area 100 nm-thick Ag film. For the cells with MoO₃, the Ag cover layer was deposited without breaking the vacuum in the same deposition system. All the Si- and SiOₓ-based layers were deposited by plasma-enhanced chemical vapor deposition (PECVD) at standard frequency or a very high frequency regime at 40 MHz (the first for intrinsic and n-type films, and the second for the p-type layers) in an MVSystems Inc. cluster tool, using previously optimized parameters [37,38].

2.3. Device Characterization

The solar cells were characterized by measuring the current–voltage characteristics under AM1.5 illumination with class AAA solar simulators. For the SHJ cells, the simulated sunlight was provided by a dual-lamp WACOM simulator, and the measurements were carried out at standard condition, after adjustment of the total irradiance based on a c-Si reference cell. The PSCs were measured in a nitrogen-filled glove box (O₂ and H₂O < 1 ppm) at 25 °C with a Keithley 2400 source measure unit (Keithley Instruments Inc., Cleveland, OH, USA). Simulated AM1.5 illumination was provided by a Photo Emission Tech solar simulator (model CT100AAA), equipped with a 150 W Xenon lamp, whose intensity was calibrated using a c-Si reference cell with a KG5 filter for 1 sunlight intensity of 100 mW/cm². Dark I-V curves were also measured with a Keithley 2400 source measure unit. The external quantum efficiency (EQE) spectra were measured with a Bentham PVE300 apparatus calibrated with an Si detector.

3. Results and Discussion

3.1. Properties of MoO₃ Films

MoO₃ layers deposited on glass were characterized with spectroscopic ellipsometry (SE) for evaluating optical properties and thickness. Due to the susceptibility of MoO₃ to air exposure, the measurements were carried out immediately after the deposition. The SE data were analyzed by modeling the samples as glass covered with a layer mathematically described with a Tauc–Lorentz dispersion dependence, with the parameters of the Tauc–Lorentz formula and the film thickness as fit parameters. With the feedback on the thickness values, a regime with a growth rate of 1 Å/s was selected, which makes the target thickness for application into solar cells (≤10 nm) easily manageable. The optical constants (refractive index, n, and extinction coefficient, k) obtained by fitting the SE data
to the model are shown in Figure 1 together with the spectra of the films used with the same function in the solar cells discussed in this work, which were likewise deposited on glass and characterized by SE. The comparison highlights the marked difference in the optical properties, with MoO\textsubscript{x} offering the advantage of negligible light absorption (parasitic absorption loss when the layer is used in the solar cell) accompanied by a low refractive index (≤2 above approximately 500 nm).

![Figure 1. Refractive index (n) and extinction coefficient (k) spectra of the films applied as hole collectors in the present work, as evaluated through spectroscopic ellipsometry.](image)

Another relevant property for application into solar cells is the work function that plays a role in the energy band alignment/misalignment within the device. The WF of 5.95 eV, measured on as-deposited samples, decreased to 5.52 eV after 24 h in air [34], confirming the strong effects of air exposure on the material properties. Such values make the material a good candidate for hole collection in solar cells.

3.2. Polymer Solar Cells

In order to test the performance of the thermally evaporated MoO\textsubscript{x} layer as an HTL in inverted PSCs, we have realized devices with the following architecture: glass/ITO/ZnO/PTB7-Th:[70]PCBM/HTL/Ag, where the HTL is MoO\textsubscript{x} or PEDOT:PSS. The investigated structure is depicted in Figure 2 together with the chemical structure of the materials used to prepare the photoactive layer. PTB7-Th is a low-band gap polymer recently used as a donor material in PSCs, showing PCE larger than 10% using an inverted device structure [39]. This material has several advantages but also some important drawbacks. On the negative side, solar cells based on PTB7-like and fullerene are generally known to be thermally unstable due to deterioration of the blend morphology at high temperatures [40,41]. On the other hand, it has been demonstrated that this type of polymer can be considered a promising material thanks to the possibility of processing it under ambient conditions using roll-to-roll or compatible production methods [42].

The current–voltage characteristics (in dark and under illuminated conditions) of the two investigated devices are reported in Figure 3. Both the dark J-V curves show a rectification behavior at ±1 V, and the reverse current is between 10\textsuperscript{-5} and 10\textsuperscript{-4} A. When comparing the two types of cells, the MoO\textsubscript{x}-based device shows a lower series resistance and a higher shunt resistance.

The values of PCE, open circuit voltage (V\textsubscript{OC}), short circuit current density (J\textsubscript{SC}), and fill factor (FF) extracted from the illuminated J-V curve (Figure 3b) are summarized in Table 1. A minimum of 5 cells was realized for each device configuration, and good reproducibility was found (within 2%) for all the extracted J-V light parameters. The best performance is obtained with the device realized by using MoO\textsubscript{x} as an HTL, which reaches a PCE of 8.67% compared to 6.23% of the PEDOT:PSS-based cell. The higher PCE value is mainly due to the improved FF (67.8% versus 58.3%) and to the higher J\textsubscript{sc} (15.9 versus 13.1 mA/cm\textsuperscript{2}).
1. When light enters the solar cells through the front contact (ITO), it is partially absorbed by the active layer, with a higher current output observed from the corresponding illuminated device. The response of the MoOₓ-based device is larger compared to the PEDOT:PSS one, confirming the improved current density.

2. The EQE of PSCs realized with different HTLs (PEDOT:PSS or MoOₓ) is reported in Figure 4. Both the devices show a broad response range from 300 to 850 nm (in agreement with the EQE reported in the literature for cells based on the same photoactive layer [39]), with values above 50% (PEDOT:PSS-based device) and 60% (MoOₓ-based device) in the range 500–700 nm. The response of the MoOₓ-based device is larger compared to the PEDOT:PSS one, confirming the improved current output observed from the corresponding illuminated J–V curves. The higher current of the device with MoOₓ as HTL is mainly due to the higher light absorption achieved by the photoactive layer in this configuration, thanks to the lower extinction coefficient of MoOₓ versus PEDOT:PSS for wavelengths above 400 nm, as found through spectroscopic ellipsometry (Figure 1). This is explained by the fact that when light enters the solar cells through the front contact (ITO), it is partially absorbed by the
active layer, while a fraction reaches the Ag back reflecting contact. Thanks to the reduced parasitic absorption losses in the MoO$_x$ layer, more light is reflected backwards into the photoactive layer, which thus has a better “second chance” to absorb it and produce a higher current output. In addition, also an improved charge collection likely contributes to enhance the EQE curve of the MoO$_x$-based device overall, which is in agreement with the higher FF.

![Figure 4](image-url)  
**Figure 4.** External quantum efficiency (solid lines, left axis) of inverted PSCs with different HTLs (PEDOT:PSS or MoO$_x$) based on the structure illustrated in Figure 2. The integrated $J_{SC}$ calculated from the external quantum efficiency (EQE) spectra is also reported (dashed lines, right axis).

### 3.3. Silicon Heterojunction Solar Cells

For the application of MoO$_x$ at the back side of SHJ solar cells, we have considered the three proof-of-concept planar structures in Figure 5, which make use of a single-layer rear-reflector (the Ag back electrode) and the same front layer stack featuring a highly transparent n-type layer [37,43]. The structure in panel (a), which uses a conventional p-type doped Si-based layer (amorphous silicon in the present case), is the reference solar cell. In (b), we have considered the replacement of the conventional hole collector with MoO$_x$, while keeping the same wafer type. We have additionally tested the (c) structure, which is the same as (b) but based on the n-type wafer. Although MoO$_x$ serves as a hole collector in both the (b) and (c) structures, there are differences in the hole transport mechanism toward the rear electrode, since holes are majority (minority) carriers in p-type (n-type) wafers.

Ultrathin MoO$_x$ is required in the devices to avoid series resistance losses due to the poor conductivity of the material (fully stoichiometric MoO$_x$ is insulating). To establish the appropriate values, we have conducted a preliminary experiment on structure (b) fabricating solar cells with variable MoO$_x$ thickness (5, 10, and 15 nm) and nominally identical intrinsic and n-type layers. Dark and illuminated current–voltage curves are shown in Figure 6 for two representative cells for each MoO$_x$ thickness value, which are selected to give direct evidence of the variability within each set of cells. Functioning solar cells have been obtained in all the cases. All the dark I-V curves show low reverse current ($<10^{-5}$ A) and proper rectification behavior (ratio of forward over reverse current $\geq 10^4$ at $\pm 1$ V), giving indication of a good passivation level and low shunt effects. No effects on the series resistance are visible when changing the MoO$_x$ thickness. The only relevant difference is the bigger variability in the leakage current among the cells with the thinnest MoO$_x$ layer. In addition, from the illuminated I-V curves, no major variation on the thickness is found for all the photovoltaic parameters ($V_{OC}$ approximately 650 mV, FF approximately 73%, $J_{SC}$ approximately 33 mA/cm$^2$), thus leaving relaxed margins for the application.
Figure 5. Schematic drawings (not to scale) of the silicon heterojunction solar cell structures investigated in this work, all with hole collectors at the back side (opposite to illumination): (a) p-type Si wafer with front emitter and conventional hole collector; (b) p-type Si wafer with front emitter and novel hole collector; (c) n-type Si wafer with novel hole collector functioning also as a rear emitter.

Figure 6. Current–voltage characteristics in the dark (a) and under AM1.5 illumination (b) of p-type based silicon heterojunction (SHJ) solar cells with an MoO$_x$ hole collector (structure (b) in Figure 5) for three MoO$_x$ thickness values.

Based on the preliminary test, the intermediate value of 10 nm has been selected for the MoO$_x$ layer thickness to be applied for the comparison of the structures in Figure 5. While in the front side application MoO$_x$ can supply a relevant optical gain thanks to the inferior parasitic absorption of light in the ultraviolet and visible region of the solar spectrum (Figure 1), here, considering that only the near infrared light reaches the bottom of the Si wafer, MoO$_x$ and (p)a-Si:H are roughly equivalent, being both transparent in the long wavelength range. Therefore, no optical gain is here pursued, but there are margins for electrical improvement together with the technological simplification related to the use of an evaporated dopant-free layer. The $J$--$V$ curves under illumination and the corresponding photovoltaic parameters for the best cells with the three designs, with nominally identical intrinsic and n-type layers, are reported in Figure 7 and Table 2, respectively. Compared to the reference cell, a considerable increase of $V_{OC}$ (38 mV) is obtained when moving from (p)a-Si:H to the MoO$_x$ collector. At the same time, FF and $J_{SC}$ remain comparable, leading to an overall increase of the conversion efficiency of approximately 1% absolute. Remarkably, the solar cells based on the n-type wafer, where the MoO$_x$ layer functions as a non-conventional emitter placed on the back side (structure (c)), show very similar characteristics and thus similar improvement against the selected reference cell. In this case, there is a margin for significant additional gain with optimization of the front stack, which here was left unchanged for comparison.
We have demonstrated the effectiveness of the MoO\textsubscript{x} hole collector at the rear of SHJ solar cells, irrespective of the wafer doping type. The significant \(V_{\text{OC}}\) increase against the reference cell is likely the effect of the improved surface passivation achieved for the Si wafer in the presence of the full (i)a-Si:H plus hole-collector stack when MoO\textsubscript{x} replaces (p)a-Si:H. It is known that when (p)a-Si:H is deposited over the ultrathin (i)a-Si:H passivation layer, the injection-level dependent lifetime drops \([38,44]\). However, the suggested more favorable condition achieved with MoO\textsubscript{x} would be challenging to demonstrate with lifetime measurements carried out on unfinished cell precursors (i.e., with unprotected MoO\textsubscript{x}), due to the high sensitivity of MoO\textsubscript{x} to air exposure.

4. Conclusions

We have reported on the application of ultrathin evaporated MoO\textsubscript{x} as a hole collector at the back side of solar cells of two very different photovoltaic technologies. In case of polymer solar cells, MoO\textsubscript{x} has been tested against a standard HTL made of PEDOT:PSS in devices with inverted architecture employing the PTB7-Th:[70]PCBM blend as a photoactive layer. The MoO\textsubscript{x}-based PSCs have shown superior performance in terms of the current output and fill factor compared to the reference with PEDOT:PSS. The superior transparency of the MoO\textsubscript{x} film allows an enhanced back reflected light inside the blend film, producing a higher photogeneration, as found from the illuminated \(J-V\) and EQE curves. The higher FF of the MoO\textsubscript{x}-based device also suggests a better charge collection efficiency.

As for the SHJ technology, we have reported on the use of MoO\textsubscript{x} as a rear dopant-free hole collector. We have demonstrated the efficacy of the approach when using both p-type and n-type Si wafers and evidenced a potential efficiency gain. In the present comparison over planar test structures with an Ag back reflecting electrode, the increase has been approximately 1\% absolute against a standard p-type amorphous silicon hole collector. The gain is linked to the increased \(V_{\text{OC}}\), which is likely due to the reduced recombination at the Si wafer surface when passivated by the (i)a-Si:H/MoO\textsubscript{x} stack in place of (i)a-Si:H/(p)a-Si:H. We have elsewhere shown that the detrimental effects of the highly defective (p)a-Si:H collection layer can be partly overcome by using p-type doped nanocrystalline Si layers with

![Figure 7. J-V characteristics under illumination of SHJ solar cells with different hole collectors ((p)a-Si:H or MoO\textsubscript{x}) based on p-type and n-type Si wafers designed as illustrated in Figure 5.](image-url)
superior electrical properties [38], but the replacement of doped Si-based layers with MoOx offers the advantage of a low-temperature low-cost fabrication technique and no need of toxic doping gases in the fabrication process. In addition, less variability is expected in the solar cell results by avoiding the delicate tuning of the deposition parameters of the doped amorphous and nanocrystalline Si layers required for obtaining the optimal optoelectronic properties of the layers.

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