Abstract. We investigate the generalized second-order Arnoldi (GSOAR) method, a generalization of the SOAR method proposed by Bai and Su [SIAM J. Matrix Anal. Appl., 26 (2005): 640–659.], and the Refined GSOAR (RGSOAR) method for the quadratic eigenvalue problem (QEP). The two methods use the GSOAR procedure to generate an orthonormal basis of a given generalized second-order Krylov subspace, and with such basis they project the QEP onto the subspace and compute the Ritz pairs and the refined Ritz pairs, respectively. We develop implicitly restarted GSOAR and RGSOAR algorithms, in which we propose certain exact and refined shifts for respective use within the two algorithms. Numerical experiments on real-world problems illustrate the efficiency of the restarted algorithms and the superiority of the restarted RGSOAR to the restarted GSOAR. The experiments also demonstrate that both IGSOAR and IRGSOAR generally perform much better than the implicitly restarted Arnoldi method applied to the corresponding linearization problems, in terms of the accuracy and the computational efficiency.

1. Introduction

Consider the large QEP

\[ Q(\lambda)x = (\lambda^2 M + \lambda C + K)x = 0 \]  

with \( \|x\| = 1 \), where \( M, C, K \) are \( n \times n \) matrices with \( M \) nonsingular and \( \| \cdot \| \) is the 2-norm of a vector or matrix. Such QEP arises in a wide variety of scientific and engineering applications [4, 25]. One is often interested in a few largest eigenvalues...
in magnitude or a few eigenvalues nearest to a target \( \sigma \) in the complex plane. One of
the commonly used approaches is to linearize the QEP and then solve the linearized
problem. There are a number of linearizations available [25], of which a commonly
used one is to transform (1) to the generalized eigenvalue problem

\[
\begin{bmatrix}
-C & -K \\
I & 0
\end{bmatrix}
\begin{bmatrix}
\lambda x \\
x
\end{bmatrix}
= \lambda
\begin{bmatrix}
M & 0 \\
0 & I
\end{bmatrix}
\begin{bmatrix}
\lambda x \\
x
\end{bmatrix},
\]

which is equivalent to the standard linear eigenvalue problem

\[
\begin{bmatrix}
A & B \\
I & 0
\end{bmatrix}
\begin{bmatrix}
\lambda x \\
x
\end{bmatrix}
= \lambda
\begin{bmatrix}
\lambda x \\
x
\end{bmatrix},
\]

where \( A = -M^{-1}C, B = -M^{-1}K \). Clearly, (3) corresponds to the monic QEP

\[
(\lambda^2I - \lambda A - B)x = 0.
\]

The mathematical theory on (2) and (3) has been well established and a number
of numerical methods have been available for solving them [1, 6, 22, 24, 26]. One
of the drawbacks via linearizations is that general numerical methods do not take
the structures of (2) and (3) into account, making computations expensive and the
approximate eigenpairs possibly lose their physical structures.

To improve the computational efficiency of the Arnoldi method that is directly
applied to some linearization problem of QEP (1), Meerbergen [19] proposes a quadratic
Arnoldi (Q-Arnoldi) method, which exploits the structure of the linearization problem
to reduce the memory requirements by about a half and can compute a partial Schur
form of the linearization problem with respect to the structure of the Schur vectors. He
shows that the Q-Arnoldi method can be implicitly restarted. Some similar methods
have proposed very recently in [20, 28]. All of them have are special Arnoldi methods
applied to certain linearization problems of QEP (1), that is, each of them projects
the corresponding linearization problem rather than QEP (1) or (4) onto some Krylov
subspace whose orthonormal basis is generated efficiently by a special Arnoldi process.
For these methods, the implicit restarting technique [23] is easily applied.

In this paper, we are interested in projection methods that work on QEP (1) directly
other than its linearizations, and such methods preserve some important structures of
it. The second-order Arnoldi (SOAR) method proposed by Bai and Su [2] falls into
this category and is a Rayleigh–Ritz method. They propose a SOAR procedure that
computes an orthonormal basis of a second-order Krylov subspace generated by the
matrices \( A \) and \( B \) simultaneously. The SOAR method then projects (1) onto this
subspace and computes the Ritz pairs to approximate the desired eigenpairs of (1). A
unified and general convergence theory has recently been established in [9] for the
Rayleigh–Ritz method and the refined Rayleigh–Ritz method for the QEP, generalizing
some of the known results on the Rayleigh–Ritz method for the linear eigenvalue problem [24]. It is proved in [9] that for a sequence of projection subspaces containing increasingly accurate approximations to a desired eigenvector there is a Ritz value that converges to the desired eigenvalue unconditionally while the corresponding Ritz vector converges conditionally and may fail to converge. Alternatively, we can compute a refined Ritz vector whose unconditional convergence is guaranteed.

In the spirit of the Hessenberg-triangular decomposition of a matrix pencil, which reduces to the Hessenberg decomposition of a single matrix, Huang et al. [10] propose a semiorthogonal generalized Arnoldi (SGA) procedure for the matrix pencil resulting from some linearization of the QEP. The SGA method first generates an SGA decomposition and then computes the Rayleigh–Ritz approximations of QEP (1) with respect to the subspace defined by an orthonormal basis generated by the SGA decomposition. To overcome the possible non-convergence of Ritz vectors obtained by the SGA method, they apply the refined projection principle [11] to propose a refined SGA (RSGA) method that computes better refined Ritz vectors. On the basis of implicitly shifted QZ iterations, they have developed the implicitly restarted SGA and RSGA algorithms, abbreviated as IRSGA and IRRSGA, with certain exact shifts and refined shifts suggested, respectively.

One disadvantage of SOAR is that the implicit restarting technique is not directly applicable. In order to make implicit restarting applicable, Otto [21] proposes a modified SOAR procedure that replaces the original special starting vector by a general one. Under the assumption that there is no deflation in the modified SOAR procedure, implicit restarting is directly adapted to this procedure. However, it is hard to interpret and understand the modified SOAR method. This is unlike the SOAR method, whose convergence is related to the Arnoldi method for the linear eigenvalue problem. Wei et al. [3, 29] make a similar modification and propose a generalized second-order Arnoldi (GSOAR) method and their refined variants, for solving the QEP and higher degree polynomial eigenvalue problems. Based on the explicit restarting scheme of the Arnoldi algorithm for the linear eigenvalue problem, Wei et al. [3, 29] have developed explicitly restarted generalized Krylov subspace algorithms. Deflation and breakdown may take place in the SOAR and modified SOAR procedures, but they have completely different consequences [2, 21], where it is proved that the SOAR method will find some exact eigenpairs of QEP (1) if breakdown occurs but no eigenpair is found generally when deflation takes place. A remedy strategy is given in [2] to treat the deflation so as to continue the SOAR procedure. Similarly, deflation may occur in the GSOAR procedure, but it is not mentioned in [3, 29].

Similar to the modified SOAR procedure, implicit restarting is directly adapted to the GSOAR procedure, but it is useable only conditionally and requires that no deflation occur in implicit restarts. Once deflation takes place, implicit restarting fails completely. Therefore, one must cure deflations, so that implicit restarting can be
applied unconditionally. For the success and overall performance of implicitly restarted GSOAR type algorithms, just as the mechanism for those implicitly restarted Krylov subspace algorithms for the linear eigenvalue problem and SVD problems [8, 12, 13, 15, 16, 18], it turns out that a proper selection of the shifts is crucial. Otto [21] has proposed certain exact shifts for his implicitly restarted modified SOAR algorithm, but they could cause convergence problems since some important aspects on the QEP are ignored when determining the shifts.

In this paper, we are concerned with the GSOAR and RGSOAR methods and their implicit restarting. We will explore more properties and features of them, and consider the efficient and reliable computation of refined Ritz vectors. Particularly, we show that there is a close relationship between the subspace generated by the GSOAR procedure and a standard Krylov subspace. With help of this result, we can interpret the convergence of the GSOAR type methods. Our main concern is a reasonable selection of the shifts when implicitly restarting the GSOAR and RGSOAR algorithms. We advance certain exact shifts, different from those in [21], and refined shifts for respective use within the implicitly restarted GSOAR and RGSOAR algorithms. The refined shifts are based on the refined Ritz vectors and theoretically better than the exact shifts. Unlike the implicitly restarted algorithms for the linear eigenvalue problem, both exact and refined shift candidates are now more than the shifts allowed. We show how to reasonably select the desired shifts among them. We present an efficient algorithm to compute the exact and refined shift candidates reliably. In addition, we propose an effective approach to cure deflations in implicit restarts, so that implicit restarting is useable unconditionally.

The rest of this paper is organized as follows. In Section 2 we review the SOAR and GSOAR procedures, present some properties of them, and describe the SOAR and GSOAR methods. In Section 3, we describe the RGSOAR method and discuss some practical issues of it. In Section 4, we develop implicitly restarted GSOAR and RGSOAR algorithms with the exact and refined shifts suggested. We present an effective approach to treat deflations in implicit restarts. In Section 5, we report numerical experiments to illustrate the efficiency of the restarted algorithms and the superiority of the refined algorithm. We also compare our algorithms with IRSGA and IRRSGA [10], demonstrating that ours perform better. More importantly, we compare our algorithms with the Matlab function `eigs`, the implicitly restarted Arnoldi method applied to a commonly used linearization problem, showing that ours generally have sharp superiority to `eigs` in terms of the accuracy and the computational efficiency. Finally, we conclude the paper in Section 6.

Throughout the paper, we denote by $\| \cdot \|$ the spectral norm of a matrix and the 2-norm of a vector, by $I$ the identity matrix with the order clear from the context, by the superscripts $T$ and $^*$ the transpose and conjugate transpose of a vector or matrix, by $\mathbb{C}^k$ the complex vector space of dimension $k$ and by $\mathbb{C}^{(k+1)\times k}$ the set of $(k+1) \times k$
matrices. We denote by $\sigma_{\min}(F)$ the smallest singular value of a matrix $F$ and by the Matlab notation $A(i : j, k : l)$ the submatrix consisting of rows $i$ to $j$ and columns $k$ to $l$ of $A$.

2. THE SOAR AND GSOAR METHODS

Bai and Su [2] introduce the following concepts.

Definition 1. Let $A, B$ be matrices of order $n$ and for the vector $u \neq 0$, and define

$$r_0 = u,$$
$$r_1 = Ar_0,$$
$$r_j = Ar_{j-1} + Br_{j-2} \quad \text{for } j \geq 2.$$  

Then $r_0, r_1, r_2, \ldots, r_{k-1}$ is called a second-order Krylov sequence based on $A, B$ and $u$, and $G_k(A, B; u) = \text{span}\{r_0, r_1, r_2, \ldots, r_{k-1}\}$ a $k$-th second-order Krylov subspace.

Note that (3) is a linearization of (4). Define the matrix

$$H = \begin{bmatrix} A & B \\ I & 0 \end{bmatrix}$$

of order $2n$. For a $2n$-dimensional starting vector $v$, we can generate a Krylov subspace $K_k(H, v) = \text{span}\{v, Hv, H^2v, \ldots, H^{k-1}v\}$. Particularly, if we choose $v = [u^T, 0]^T$, we have

$$\begin{bmatrix} r_j \\ r_{j-1} \end{bmatrix} = H^j v, \quad j \geq 0 \text{ with } r_{-1} = 0.$$  

We observe the fundamental relation

$$K_k(H, v) \subseteq G_k^2(A, B; u),$$

where $G_k^2(A, B; u)$ is the subspace generated by the vector set

$$\left\{ \begin{bmatrix} r_0 \\ 0 \end{bmatrix}, \begin{bmatrix} r_1 \\ 0 \end{bmatrix}, \ldots, \begin{bmatrix} r_{k-1} \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ r_0 \end{bmatrix}, \begin{bmatrix} 0 \\ r_1 \end{bmatrix}, \ldots, \begin{bmatrix} 0 \\ r_{k-1} \end{bmatrix} \right\}.$$  

Due to the equivalence of QEP (4) and the eigenproblem of $H$, relation (7) shows that if the eigenvector $[\lambda x^T, x^T]^T$ is contained in $K_k(H, [u^T, 0]^T)$, then the eigenvector $x$ of QEP (4) is contained in $G_k(A, B; u)$. By continuity, if there is a good approximation to $[\lambda x^T, x^T]^T$ in $K_k(H, [u^T, 0]^T)$, then there must be a good approximation to $x$ contained in $G_k(A, B; u)$.

Bai and Su [2] propose the following procedure for computing an orthonormal basis $\{q_j\}_{j=1}^k$ of $G_k(A, B; u)$ and an auxiliary vector sequence $\{p_j\}$ generating $G_{k-1}(A, B; u)$.
Algorithm 1. SOAR procedure

1: \( q_1 = u / \| u \|, \ p_1 = 0 \)
2: for \( j = 1, 2, \ldots, k \) do
3: \( r = Aq_j + Bp_j \)
4: \( s = q_j \)
5: for \( i = 1, 2, \ldots, j \) do
6: \( t_{ij} = q_i^* r \)
7: \( r = r - q_i t_{ij} \)
8: \( s = s - p_i t_{ij} \)
9: end for
10: \( t_{j+1j} = \| r \| \)
11: if \( t_{j+1j} = 0 \), stop
12: \( q_{j+1} = r / t_{j+1j} \)
13: \( p_{j+1} = s / t_{j+1j} \)
14: end for

The following basic results hold for this algorithm; see [2].

**Theorem 1.** Define \( Q_k = [q_1, q_2, \ldots, q_k] \) and \( P_k = [p_1, p_2, \ldots, p_k] \) and \( \hat{T}_k = \begin{bmatrix} T_k \\ t_{k+1k} e_k^* \end{bmatrix} = [t_{ij}] \in \mathbb{C}^{(k+1)\times k} \). If Algorithm 1 does not stop before step \( k \), then we have

\[
\text{span}\{Q_k\} = G_k(A, B; u)
\]

and the \( k \)-step SOAR decomposition

\[
H \begin{bmatrix} Q_k \\ P_k \end{bmatrix} = \begin{bmatrix} Q_{k+1} \\ P_{k+1} \end{bmatrix} \hat{T}_k,
\]

where \( Q_{k+1} = [Q_k, q_{k+1}] \), \( P_{k+1} = [P_k, p_{k+1}] \).

Before proceeding, we introduce the following definition.

**Definition 2.** [2]. If \( r_i, i = 0, 1, \ldots, j \) are linearly dependent but \( \{r_i^T, r_{i-1}^T\}^T \), \( i = 0, \ldots, j \) with \( r_{-1} = 0 \) are not, we call this situation deflation; if both \( \{r_i\} \) and \( \{r_i^T, r_{i-1}^T\} \) are linearly dependent at step \( j \), we call this situation breakdown.

According to Definition 2, if Algorithm 1 stops prematurely at step \( j < k \), then either deflation or breakdown must occur at that step. Deflation means that \( \mathcal{G}_{j+1}(A, B; u) = \mathcal{G}_j(A, B; u) \) but \( \mathcal{K}_{j+1}(H, v) \neq \mathcal{K}_j(H, v) \), so the Arnoldi process on \( H \) does not terminate at step \( j \). As a result, when deflation occurs at step \( j \), \( \mathcal{K}_j(H, v) \) does not contain any exact eigenvector of \( H \), which, from (7), implies that \( \mathcal{G}_j(A, B; u) \) may not contain any exact eigenvector of QEP (1). Therefore, deflation must be remedied to continue the algorithm.

Bai and Su [2] present the following algorithm that detects and remedies deflation.
Algorithm 2. SOAR procedure with deflation remedy

1: \( q_1 = u/\|u\|, p_1 = 0 \)
2: for \( j = 1, 2, \ldots, k \) do
3: \( r = Aq_j + Bp_j \)
4: \( s = q_j \)
5: for \( i = 1, 2, \ldots, j \) do
6: \( t_{ij} = q_i^* r \)
7: \( r = r - q_i t_{ij} \)
8: \( s = s - p_i t_{ij} \)
9: end for
10: \( t_{j+1} = \|r\| \)
11: if \( t_{j+1} = 0 \) then
12: if \( s \in \text{span}\{p_i|i : q_i = 0, 1 \leq i \leq j\} \) then
13: break
14: else deflation
15: reset \( t_{j+1} = 1 \)
16: \( q_{j+1} = 0 \)
17: \( p_{j+1} = s \)
18: end if
19: else
20: \( q_{j+1} = r/t_{j+1} \)
21: \( p_{j+1} = s/t_{j+1} \)
22: end if
23: end for

In the procedure, if deflation occurs, we simply set \( t_{j+1} \) to one and take \( q_{j+1} = 0 \).

To decide if \( s \in \text{span}\{p_i|i : q_i = 0, 1 \leq i \leq j\} \), the Gram–Schmidt orthogonalization with refinement is used [2, 21]. When deflation occurs, the nonzero vectors in the sequence \( \{q_j\} \) are still orthonormal and span the second-order Krylov subspace \( G_k(A, B; u) \) with the dimension smaller than \( k \). We refer the reader to Bai and Su [2] for details.

We point out that Theorem 1 is true for Algorithm 2 but there are zero columns in \( Q_k \) when deflation occurs.

It is easily checked that a serious disadvantage of the SOAR procedure is that the implicit restarting technique is not applicable since the updated \( p_1 \) is not zero any more. Several researchers have proposed replacing \( p_1 = 0 \) in Algorithms 1–2 by a nonzero one [3, 21, 29]. This leads to the following generalized second-order Krylov sequence and subspace; see [3, 29].

**Definition 3.** Let \( A \) and \( B \) be \( n \times n \) matrices and for vectors \( u_1, u_2 \in \mathbb{C}^n \), and define

\[
\begin{align*}
r_0 &= u_1, \\
r_1 &= Ar_0 + Bu_2, \\
r_j &= Ar_{j-1} + Br_{j-2} \quad \text{for } j \geq 2.
\end{align*}
\]

Then \( r_0, r_1, r_2, \ldots, r_{k-1} \) is called a generalized second-order Krylov sequence based on \( A, B \) and \( u_1, u_2 \) and \( G_k(A, B; u_1, u_2) = \text{span}\{r_0, r_1, r_2, \ldots, r_{k-1}\} \) the \( k \)-th generalized second-order Krylov subspace.
Obviously, \( G_k(A, B; u_1, 0) = G_k(A, B; u_1) \). For a general \( \tilde{v} = [u_1^T, u_2^T]^T \), it is seen that

\[
\begin{bmatrix}
  r_j \\
  r_{j-1}
\end{bmatrix} = H^j \tilde{v}, \quad j \geq 1.
\]

But different from (7), since \( u_2 \) is a general vector, the fundamental relation now becomes

\[
HK_{k-1}(H, \tilde{v}) = \text{span}\{H\tilde{v}, \ldots, H^{k-1}\tilde{v}\} \subseteq G^2_k(A, B; u_1, u_2),
\]

where \( G^2_k(A, B; u_1, u_2) \) is the subspace generated by the vector set

\[
\left\{ \begin{bmatrix} r_0 \\ 0 \end{bmatrix}, \begin{bmatrix} r_1 \\ 0 \end{bmatrix}, \ldots, \begin{bmatrix} r_{k-1} \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ r_0 \end{bmatrix}, \begin{bmatrix} 0 \\ r_1 \end{bmatrix}, \ldots, \begin{bmatrix} 0 \\ r_{k-1} \end{bmatrix} \right\}.
\]

Note that if the eigenvector \([\lambda x^T, x^T]^T\) is contained in \( K_{k-1}(H, \tilde{v}) \) then it also lies in the subspace \( HK_{k-1}(H, \tilde{v}) \). If this is the case, (11) shows that the eigenvector \( x \) of QEP (1) is contained in \( G_k(A, B; u_1, u_2) \). More generally, by continuity, it is deduced from the above that if \( K_{k-1}(H, \tilde{v}) \) has a good approximation to \([\lambda x^T, x^T]^T\) then \( HK_{k-1}(H, \tilde{v}) \) has one too, which, in turn, means that there must be a good approximation to \( x \) contained in \( G_k(A, B; u_1, u_2) \).

Analogous to Algorithm 2, we can present a GSOAR procedure, i.e., Algorithm 3, that remedies deflation and generates the vector sequence \( \{q_j\} \), whose nonzero ones form an orthonormal basis of \( G_k(A, B; u_1, u_2) \). We point out that the GSOAR procedure in [3, 29] is the same as Algorithm 1 except that \( p_1 = 0 \) in line 1 is replaced by a general vector \( p_1 = u_2/\|u_2\| \).

**Algorithm 3. GSOAR procedure with deflation remedy**

1: \( q_1 = \frac{u_1}{\|u_1\|}, p_1 = \frac{u_2}{\|u_2\|} \).
2: for \( j = 1, 2, \ldots, k \) do
3: \( r = Aq_j + Bp_j \)
4: \( s = q_j \)
5: for \( i = 1, 2, \ldots, j \) do
6: \( t_{ij} = q_i^*r \)
7: \( r = r - t_{ij}q_i \)
8: \( s = s - t_{ij}p_i \)
9: end for
10: \( t_{j+1j} = \|r\| \)
11: if \( t_{j+1j} = 0 \)
12: if \( s \in \text{span}\{p_i | q_i = 0, 1 \leq i \leq j\} \)
13: break
14: else deflation
reset $t_{j+1} = 1$
$q_{j+1} = 0$
$p_{j+1} = s$
end if
else
$q_{j+1} = r/t_{j+1}$
$p_{j+1} = s/t_{j+1}$
end if
end for

It is direct to justify that Theorem 1 holds for this algorithm with a general $p_1 = u_2/\|u_2\|$, that is, we have

$$\text{span}\{Q_k\} = G_k(A, B; q_1, p_1)$$

with $q_1$ and $p_1$ normalized and the $k$-step GSOAR decomposition (9) if the algorithm does not break down before step $k$.

Otto [21] defines the modified second-order Krylov sequence as $r_0, r_1, r_2, \ldots, r_{k-1}$, $u_2$ and the modified second-order Krylov subspace of dimension $k + 1$ generated by the vector sequence. After the orthonormal $q_1, q_2, \ldots, q_{k+1}$ are generated, he orthonormalizes $u_2$ against them to get $q_{k+2}$. This is called the modified SOAR procedure. A disadvantage of it is that there is no compact relationship (7) or (11). So it is hard to interpret such a modified subspace and establish definitive results on breakdown and deflation.

The GSOAR method is a Rayleigh–Ritz method, and it projects the large QEP (1) onto $G_k(A, B; u_1, u_2)$ by imposing the Galerkin condition, leading to the $k$-dimensional QEP

$$(\theta^2 M_k + \theta C_k + K_k)g = 0$$

with $\|g\| = 1$, where $M_k = Q_k^*MQ_k$, $C_k = Q_k^*CQ_k$ and $K_k = Q_k^*KQ_k$. Let the $(\theta, g)$ be the eigenpairs of (12). Then the GSOAR method uses the Rayleigh–Ritz pairs $(\theta, y(= Q_kg))$ to approximate some of the eigenpairs of (1). We comment that if deflation occurs then $Q_k$ consists of only nonzero orthonormal vectors $q_j$ and the dimension of (12) is smaller than $k$.

3. A REFINED GSOAR (RGSOAR) METHOD

As is known, the Rayleigh–Ritz method may fail to converge for computing eigenvectors of the linear eigenvalue problem and the QEP; see [17] and [9], respectively. To correct this deficiency, a refined projection principle is proposed in [11] (see also [24, 26]) for the linear eigenvalue problem, which leads to the refined Rayleigh–Ritz method. The refined method extracts the best approximate eigenvectors from a given subspace in the sense that the residuals formed with certain approximate eigenvalues
available are minimized in the sense of 2-norm over the subspace. A refined GSOAR (RGSOAR) method has been proposed in [3, 29]. We next describe it and give more details on some practical issues.

Suppose that we have computed the Ritz values $\theta$ by the GSOAR method and select $m$ ones of them to approximate $m$ desired eigenvalues of (1). For each chosen $\theta$, the RGSOAR method seeks a unit length vector $\tilde{u} \in G_k(A, B; u_1, u_2)$ satisfying the optimal requirement

$$
\|(\theta^2 M + \theta C + K)\tilde{u}\| = \min_{u \in G_k(A, B; u_1, u_2)} \|(\theta^2 M + \theta C + K)u\| \tag{13}
$$

and uses it as an approximate eigenvector, called the refined Ritz vector. The pairs $(\theta, \tilde{u})$ are also called the refined Rayleigh–Ritz approximations. Since the (non-zero) columns of $Q_k$ form an orthonormal basis of $G_k(A, B; u_1, u_2)$, (13) amounts to seeking a unit length vector $\tilde{z} \in C^k$ such that $\tilde{u} = Q_k \tilde{z}$ with

$$
\tilde{z} = \arg \min_{z \in C^k} \|(\theta^2 M + \theta C + K)Q_k z\|, \tag{14}
$$

the right singular vector of the matrix $\theta^2 MQ_k + \theta C Q_k + K Q_k$ associated with its smallest singular value $\sigma_{\min}(\theta^2 MQ_k + \theta C Q_k + K Q_k)$. However, the direct computation of its SVD may be expensive. Precisely, assume that the matrix is real and $k \ll n$. Then the cost of Golub–Reinsch’s SVD algorithm is about $4nk^2$ flops, and that of Chan’s SVD algorithm is about $2nk^2$ flops [6, p. 254]. Keep in mind that $m$ is the number of the desired eigenpairs. The CPU time costs are then $4nmk^2$ and $2nmk^2$ flops, respectively.

The first author in [14] has proposed a cross-product matrix-based algorithm for computing the SVD of a matrix, which can be much more efficient than the above standard SVD algorithms. Applying the algorithm to (13), we form the cross-product matrix

$$
B_k = (\theta^2 MQ_k + \theta C Q_k + K Q_k)^* (\theta^2 MQ_k + \theta C Q_k + K Q_k),
$$

which is the Hermitian (semi-)positive definite. $\tilde{z}$ is then the eigenvector of $B_k$ associated with its smallest eigenvalue $\sigma_{\min}(\theta^2 MQ_k + \theta C Q_k + K Q_k)$. We compute the eigensystem of $B_k$ by the QR algorithm to get $\tilde{z}$. In finite precision arithmetic, the computed eigenvector is an approximation to $\tilde{z}$ with accuracy $O(\epsilon_{\text{mach}})$ provided that the second smallest singular value of $\theta^2 MQ_k + \theta C Q_k + K Q_k$ is not very close to the smallest one, where $\epsilon_{\text{mach}}$ is the machine precision.

Let us now look at the computational cost of this algorithm. Define

$$
W_1 = MQ_k, \ W_2 = CQ_k, \ W_3 = KQ_k,
$$

which are available when forming the projected QEP and do not need extra cost. Then
where the bar denotes the complex conjugate of a scalar. Assume that $W_1, W_2$ and $W_3$ are real and note that $B_k$ is Hermitian for a complex $\theta$ and real symmetric for a real $\theta$. Then we only need to form the upper (lower) triangular part of $B_k$, which involves the upper (lower) triangular parts of the nine matrices $W_i^* W_j, i, j = 1, 2, 3$. All these cost about $9nk^2$ flops. With these nine $W_i^* W_j$ available, we only need $O(k^2)$ flops to form $B_k$ for either a real or complex $\theta$, negligible to $9nk^2$ flops. So, we CPU timely need $9nk^2$ flops to form $m$ Hermitian matrices $B_k$ for $m$ approximate eigenvalues $\lambda_i$. We then compute the complete eigensystems of these $B_k$ by the QR algorithm using $O(mk^3)$ flops. Therefore, we can compute $m$ right singular vectors $\tilde{z}$ using about $9nk^2$ flops when $mk \ll n$, a natural requirement in practice. As a result, a simple comparison indicates that such cross-product based algorithm is more efficient than Golub–Reinsch’s SVD algorithm when $m \geq 3$ and Chan’s SVD algorithm when $m \geq 5$.

We can now present a basic (non-restarted) RGSOAR algorithm.

**Algorithm 4. The RGSOAR algorithm**

1. Given the starting vectors $u_1, u_2$, run the GSOAR procedure to generate an orthonormal basis $Q_k$ of $\mathcal{G}_k(A, B; u_1, u_2)$.
2. Compute $W_1 = MQ_k$, $W_2 = CQ_k$ and $W_3 = KQ_k$.
3. Compute $M_k = Q_k^* W_1$, $C_k = Q_k^* W_2$ and $K_k = Q_k^* W_3$, solve the projected QEP

$$
(\theta_i^2 M_k + \theta_i C_k + K_k)g_i = 0, \tag{16}
$$

and select $m$ Ritz values $\theta_i$ as approximations to the $m$ desired eigenvalues $\lambda_i$.

1. For each chosen $\theta_i, 1 \leq i \leq m$, form $B_k$, and compute the eigenvector $\bar{z}_i$ of $B_k$ associated with its smallest eigenvalue and the refined Ritz vector $\bar{u}_i = Q_k \bar{z}_i$.
2. Test convergence of $(\theta_i, \bar{u}_i)$ by computing the relative residual norms

$$\frac{|| (\theta_i^2 M + \theta_i C + K) \bar{u}_i ||}{|\theta_i|\| M \|_1 + |\theta_i|\| C \|_1 + \| K \|_1}, \quad i = 1, 2, \ldots, m.$$  

4. IMPLICITLY RESTARTED ALGORITHMS

This section consists of three subsections. In Section , under the assumption that no deflation occurs, we describe how to implicitly restart the GSOAR procedure. In Section , we discuss how to select best possible shifts, and propose exact and refined shifts for respective use within implicitly restarted GSOAR and RGSOAR algorithms.
In Section 4.3, we present an effective approach to cure deflation in implicit restarts, so that implicit restarting can be run unconditionally.

4.1. Implicit restarts

As step \( k \) increases, the GSOAR and RGSOAR methods become expensive and impractical due to storage requirement and/or computational cost. So restarting is generally necessary. That is, for a given maximum \( k \) and the subspace \( G_k(A, B; q_1, p_1) \) with \( q_1 \) and \( p_1 \) normalized, if the methods do not converge yet, based on the information available, we select new unit length vectors \( q_1^+ \) and \( p_1^+ \) to construct a better subspace \( G_k(A, B; q_1^+, p_1^+) \) that contains richer information on the desired eigenvectors \( x \). We then extract new better approximate eigenpairs with respect to \( G_k(A, B; q_1^+, p_1^+) \). Proceed in such a way until the methods converge.

If no deflation occurs, it is direct to adapt the implicit restarting scheme [23] to the modified SOAR procedure in [21] and the GSOAR procedure. Given \( p \) shifts \( \mu_1, \mu_2, \ldots, \mu_p \), performing \( p \) implicit shifted QR iterations on \( T_k \) yields the relation

\[
(T_k - \mu_1 I) \cdots (T_k - \mu_p I) = V_k R,
\]

where \( V_k \) is a \( k \times k \) orthogonal (unitary) matrix and \( R \) is upper triangular. Specifically, \( V_k \) has only \( p \) nonzero subdiagonals. Adapted from the derivation of implicitly restarting the standard Arnoldi process [23], we can establish the following result for the GSOAR procedure.

**Theorem 2.** Given \( p \) shifts \( \mu_1, \ldots, \mu_p \), perform \( p \) steps of implicit shifted QR iterations on \( T_k \). Let \( \psi(T_k) = V_k R_k \) with \( \psi(\mu) = \prod_{j=1}^{p} (\mu - \mu_j) \), and define \( Q_k^+ = Q_k V_k \) and \( T_k^+ = V_k^* T_k V_k \). Assume that no deflation occurs in the \( k = m + p \)-step GSOAR decomposition (9). Then we have an updated \( m \)-step GSOAR decomposition

\[
H \left[ \begin{array}{c} Q_m^+ \\ P_m^+ \end{array} \right] = \left[ \begin{array}{c} Q_m^+ \\ P_m^+ \end{array} \right] T_m^+ + \tilde{f}_{m+1}^- \left[ \begin{array}{c} q_{m+1}^- \\ p_{m+1}^- \end{array} \right] e_m^*
\]

starting with \( \left[ \begin{array}{c} q_1^+ \\ p_1^+ \end{array} \right] \), where \( Q_m^+ = Q_k V_k(:, 1 : m) \), \( P_m^+ = P_k V_k(:, 1 : m) \), \( T_m^+ = T_k^+(1 : m, 1 : m) \) is upper Hessenberg and

\[
\left[ \begin{array}{c} q_{m+1}^- \\ p_{m+1}^- \end{array} \right] = \frac{1}{f_{m+1}^-} \left[ \begin{array}{c} q_{m+1}^- \\ p_{m+1}^- \end{array} \right]
\]

\[
f_{m+1}^- = t_{m+1}^- \left[ \begin{array}{c} q_{m+1}^- \\ p_{m+1}^- \end{array} \right] + t_{k+1}^+ V_k(k, m) \left[ \begin{array}{c} q_{k+1}^- \\ p_{k+1}^- \end{array} \right]
\]

\[
\tilde{f}_{m+1}^- = \| t_{m+1}^- q_{m+1}^- + t_{k+1}^+ V_k(k, m) q_{k+1}^- \|
\]

with \( V_k(k, m) \) the entry of \( V_k \) in position \( (k, m) \).
Theorem 2 states that if no deflation occurs then we have naturally obtained an $m$-step GSOAR decomposition (17) after $p$ implicit shifted QR iterations are run on $T_k$, thus generating an orthonormal basis $\{q_j\}_{j=1}^m$ of the $m$-dimensional subspace $G_m(A, B; q_1^+, p_1^+)$. Decomposition (17) is then extended to a $k$-step one from step $m + 1$ upwards in a standard way other than from scratch, producing an orthonormal basis $\{q_j^+\}_{j=1}^k$ of the updated $k$-dimensional subspace $G_k(A, B; q_1^+, p_1^+)$. Analogous to the proof of the result on updated starting vectors in [23], it is direct to justify the following theorem.

**Theorem 3.** It holds that

$$
\begin{bmatrix}
q_1^+ \\
p_1^+
\end{bmatrix} = \frac{1}{\tau} \psi(H) \begin{bmatrix}
q_1 \\
p_1
\end{bmatrix},
$$

with $\psi(\lambda) = \prod_{j=1}^p (\lambda - \mu_j)$ and $\tau$ a normalizing factor.

### 4.2. The selection of shifts

The selection of the shifts is one of the keys for the success and overall efficiency of implicitly restarted GSOAR and RGSOAR algorithms. In this subsection we propose the corresponding best possible shifts for respective use within each algorithm.

Assume that $H$ is diagonalizable. It is shown in, e.g., [22], that if the starting vector $\tilde{v}$ is a linear combination of $m$ eigenvectors of $H$ then $K_m(H, \tilde{v})$ is an invariant subspace. Therefore, a fundamental principle of restarting is to select a better vector $\tilde{v}^+$, in some sense, from the current $K_k(H, \tilde{v})$ as an updated starting vector that amplifies the components of the desired eigenvectors and simultaneously damps those of the unwanted ones, so that the updated $K_k(H, \tilde{v}^+)$ contains more accurate approximations to the $m$ desired eigenvectors. For implicit restarting, based on formulas for updated starting vectors like (18), for the linear eigenvalue problem and the computation of a partial SVD, it has been shown in [12, 13] and [15, 16] that such goal is achieved by selecting the shifts to approximate some of the unwanted eigenvalues or singular values as best as possible within the framework of the underlying method. A general result is that the better the shifts approximate the unwanted eigenvalues, the richer information on the desired eigenvectors is contained in the updated starting vector, so that a better Krylov subspace is generated.

Motivated by the above results, we now investigate a reasonable selection of shifts for use within implicitly restarted GSOAR and RGSOAR algorithms. Observe that the projected QEP (12) of the large QEP (1) over $\text{span}\{Q_k\}$ amounts to the generalized eigenvalue problem

$$
\begin{bmatrix}
-C_k & -K_k \\
I & 0
\end{bmatrix}
\begin{bmatrix}
\theta g \\
g
\end{bmatrix} = \begin{bmatrix}
M_k & 0 \\
0 & I
\end{bmatrix}
\begin{bmatrix}
\theta g \\
g
\end{bmatrix},
$$

where $C_k = C(A, B; q_1^+, p_1^+)$ and $K_k = K(A, B; q_1^+, p_1^+)$. Theorem 3 ensures that for any $\mu_j$ with $\mu_j \neq \lambda_1$ this can be done by selecting $\tilde{v}^+ = \begin{bmatrix} q_1^+ & p_1^+ \end{bmatrix}$ as the updated starting vector, so that

$$
\begin{bmatrix}
-C_k & -K_k \\
I & 0
\end{bmatrix}
\begin{bmatrix}
\theta g \\
g
\end{bmatrix} = \begin{bmatrix}
\theta q_1 \\
\theta p_1
\end{bmatrix}. 
$$

Therefore, the selection of the shifts is reduced to that of the unwanted eigenvalues $\mu_j$ with $\mu_j \neq \lambda_1$. The next subsection shows how to do that.
which is the projected problem of large generalized eigenvalue problem (2) over the subspace $G_2^k(A, B; u_1, u_2)$ (c.f. (11)) spanned by the (nonzero) columns of

$$
\hat{Q}_{2k} = \begin{bmatrix}
Q_k & 0 \\
0 & Q_k
\end{bmatrix}.
$$

The above problem amounts to the standard linear eigenvalue problem

$$
\begin{bmatrix}
-M_k^{-1}C_k & -M_k^{-1}K_k \\
I & 0
\end{bmatrix}
\begin{bmatrix}
\theta g \\
g
\end{bmatrix} = \theta
\begin{bmatrix}
\theta g \\
g
\end{bmatrix}.
$$

(18) indicates that we should select the shifts $\mu_j, j = 1, 2, \ldots, m$ as the best possible approximations to the unwanted eigenvalues of $H$ so as to generate increasingly better updated subspaces $K_k(H, v^+)$ and $HK_{k-1}(H, v)$ with $v^+ = [q_1^T, p_1^T]^T$. In terms of (11) and the comments followed, this, in turn, leads to increasingly better updated $G_2^k(A, B; q_1^+, p_1^+)$ that contains increasingly better approximations to the $m$ desired eigenvectors of $H$. As a result, $G_k(A, B; q_1^+, p_1^+)$ contains more accurate approximations to the desired eigenvectors of (1). So, just as for the linear eigenvalue problem, we should choose shifts for each implicitly restarted GSOAR type algorithm in the sense that they are best possible approximations to some of the unwanted eigenvalues of (1).

For the Rayleigh–Ritz method with respect to a given subspace, the Ritz values can be considered as the best approximations available to some eigenvalues of (1). Otto [21] proposed exact second-order shifts for his implicitly restarted modified SOAR algorithm. Adapted here, one solves the projected QEP (12) and selects $m$ Ritz values $\theta_i$ as approximations to the desired eigenvalues. Then the unwanted Ritz values are shift candidates, called the exact second-order shift candidates. A problem is that there are $2k - m$ shift candidates, while for (17) the number $p$ of shifts must not exceed $k - m$. One must select $p = k - m$ shifts among the $2k - m$ candidates. Otto simply suggested to take any $p = k - m$ shifts among $2k - m$ ones. We should point out that this situation is unlike implicitly restarted Arnoldi type algorithms for the linear eigenvalue problem, where the the maximum number of shifts is just that of candidates; see [23] and [12, 13, 15, 16].

However, the above selection of exact second-order shifts is problematic and susceptible to failure, as elaborated below. It is crucial to keep in mind a basic fact that the QEP may often have two distinct eigenvalues that share the same eigenvector [25]. This means that, for QEP (12), some of the shift candidates and some of the $m$ Ritz values used to approximate the desired eigenvalues may share common eigenvector(s). Therefore, if it is unfortunate to take such candidates for shifts, restarting will filter out the information on the corresponding desired eigenvectors and thus makes implicitly restarted GSOAR algorithms perform poorly.

In order to avoid the above deficiency, we propose new shift candidates for the implicitly restarted GSOAR and RGSOAR algorithms, respectively, and show how
to reasonably select the shifts among the candidates. We first consider the GSOAR method. Project QEP (1) onto the orthogonal complement of $\text{span}\{y_1, \ldots, y_m\}$ with respect to $G_k(A, B; q_1, p_1)$, where $y_1, \ldots, y_m$ are the Ritz vectors approximating the desired eigenvectors $x_1, \ldots, x_m$. Then we obtain a $p$-dimensional projected QEP and compute its $2p$ eigenvalues. A remarkable consequence is that these $2p$ eigenvalues must be approximations to some of the unwanted eigenvalues of QEP (1) because the information on $x_1, \ldots, x_m$ has been removed from $G_k(A, B; q_1, p_1)$. So we can use any $p$ ones of these $2p$ candidates as shifts. To be unique, we choose the $p$ ones farthest from the Ritz values $\theta_i, i = 1, 2, \ldots, m$ that are used to approximate the desired eigenvalues $\lambda_1, \ldots, \lambda_m$. The motivation of this choice is that, based on (18), these shifts can be better to amplify the information of $\tilde{u}^+$ on the desired eigenvectors and dampen the components of undesired eigenvectors in $\tilde{v}^+$.

If we are interested in the $m$ eigenvalues nearest to a target $\sigma$ and/or the associated eigenvectors, QEP (1) can be equivalently transformed to a shift-invert QEP; see the end of this subsection. In this case, we select the $p$ Ritz values among $2p$ candidates farthest from $\sigma$ as shifts. Such selection of shifts is motivated by an idea from [15, 16], where some of the shifts are taken to be unwanted Ritz values farthest from the wanted approximate singular values. It was argued there that this selection can better dampen those components of the unwanted singular vectors and meanwhile amplify the components of the desired singular vectors.

We now turn to the selection of shifts for the RGSOAR algorithm. Algorithm 4 computes the refined Ritz vectors $\tilde{u}_i$, which are generally more and can be much more accurate than the Ritz vectors $y_i$ [9, 17]. The first author [12, 13] has proposed certain refined shifts for the refined Arnoldi method and the refined harmonic Arnoldi method for the linear eigenvalue problem. It is shown that the refined shifts are generally better than the corresponding exact shifts and can be computed efficiently and reliably. In the same spirit, we next propose certain refined shifts for the RGSOAR algorithm.

Since the refined Ritz vectors $\tilde{u}_i, i = 1, 2, \ldots, m$ are more accurate than the corresponding $y_i$, the orthogonal complement of $\text{span}\{\tilde{u}_1, \ldots, \tilde{u}_m\}$ with respect to $G_k(A, B; q_1, p_1)$ contains richer information on the unwanted eigenvectors than the orthogonal complement of $\text{span}\{y_1, \ldots, y_m\}$ with respect to $G_k(A, B; q_1, p_1)$. As a result, the eigenvalues of the projected QEP of QEP (1) onto this orthogonal complement are more accurate approximate eigenvalues than the exact shift candidates described above. We call them refined shift candidates. We use the same approach as above to select $p$ ones among them as shifts, called the refined shifts, for use within the implicitly restarted RGSOAR algorithm.

Finally, we show how to compute the exact and refined shifts efficiently and reliably. We take the refined shifts as example. The computation of exact shifts is analogous. Recall $\tilde{u}_i = Q_k \tilde{z}_i, i = 1, 2, \ldots, m$, and write $\tilde{Z}_m = [\tilde{z}_1, \ldots, \tilde{z}_m]$. If QEP (1) is real and two columns $\tilde{z}_i$ and $\tilde{z}_{i+1}$ of $\tilde{Z}_m$ are complex conjugate, we replace them by their
normalized real and imaginary parts, respectively, so that the resulting $Z_m$ is real. We then make the full QR decomposition

$$Z_m = [U_m, U_\perp] \begin{bmatrix} R_m \\ 0 \end{bmatrix},$$

where $U_m$ and $U_\perp$ are $k \times m$ and $k \times p$ column orthonormal matrices, respectively, and $R_m$ is $m \times m$ upper triangular. We use the Matlab built-in function qr.m to compute the decomposition in experiments. This costs $O(k^3)$ flops, negligible to the cost of the $k$-step GSOAR procedure. Obviously, it holds that

$$\text{span}\{\tilde{u}_1, \ldots, \tilde{u}_m\} = \text{span}\{Q_k U_m\}, \quad \text{span}\{[Q_k U_m, Q_k U_\perp]\} = G_k(A, B; q_1, p_1).$$

Therefore, $Q_k U_\perp$ is an orthonormal basis of the orthogonal complement of $\text{span}\{\tilde{u}_1, \ldots, \tilde{u}_m\}$ with respect to $G_k(A, B; q_1, p_1)$. It is direct to justify that the projected QEP of the original QEP (1) onto $\text{span}\{Q_k U_\perp\}$ is just the projected QEP of the small QEP (16) onto $\text{span}\{U_\perp\}$. So, we form the projected QEP of the original QEP (1) onto $\text{span}\{Q_k U_\perp\}$ at cost of $O(k^3)$ flops. We then compute its $2p$ eigenvalues using $O(p^3)$ flops and select $p$ ones among them as the refined shifts. Since $p < k$, the CPU time cost of computing the refined shifts is $O(k^3)$ flops. For the exact shifts, recall the Ritz vectors $y_i = Q_k g_i, i = 1, 2, \ldots, m$. Write $G_m = [g_1, \ldots, g_m]$ and replace $Z_m$ by it. We then compute the exact shifts in the same way as above.

Having done the above, we have finally developed the following Algorithm 5.

**Algorithm 5. The implicitly restarted GSOAR type algorithms**

1. Given unit length starting vectors $q_1$ and $p_1$, the number $m$ of desired eigenpairs and the number $p$ of shifts $p$ satisfying $p \leq m - k$, run the $k$-step GSOAR procedure to generate $Q_k$.

2. Do until convergence

   Project QEP (1) onto $\text{span}\{Q_k\}$ to get QEP (12), select $m$ Ritz pairs $(\theta_i, y_i)$ or refined Ritz pairs $(\theta_i, \tilde{u}_i)$ as approximations to the $m$ desired eigenpairs, respectively, and determine their convergence.

3. If not converged, compute the $p$ exact shifts or refined shifts, and implicitly restart the GSOAR method or the RGSOAR method, respectively.

4. EndDo

Algorithm 5 includes two algorithms: the implicitly restarted GSOAR algorithm with the exact shifts and RGSOAR algorithm with the refined shifts, abbreviated as IGSOAR and IRGSOAR here and hereafter. They can be used to compute a number of largest eigenvalues in magnitude and the associated eigenvectors of QEP (1). We determine the convergence of a Ritz pair $(\theta, y)$ by requiring

$$\frac{\|\theta^2 M + \theta C + K\| y\|}{|\theta|^2 M\|_1 + |\theta| C\|_1 + K\|_1} \leq \text{tol},$$

(20)
where \( tol \) is a user-prescribed accuracy. For the convergence of a refined Ritz pair \((\theta, \tilde{u})\), we replace the above \( y \) by \( \tilde{u} \).

If the \( m \) eigenvalues closest to a given target \( \sigma \) are desired, we use the shift-invert transformation \( \rho = \frac{1}{\lambda - \sigma} \) with \( \det(Q(\sigma)) \neq 0 \) to transform QEP (1) to the new QEP

\[
Q_\sigma(\rho)x = (\rho^2 M_\sigma + \rho C_\sigma + K_\sigma)x = 0,
\]

where \( M_\sigma = \sigma^2 M + \sigma C + K \) is nonsingular as \( \det(M_\sigma) = \det(Q(\sigma)) \neq 0 \), \( C_\sigma = C + 2\sigma M \), \( K_\sigma = M \). We then apply the previous analysis and algorithms to (21). Let \((\tilde{\rho}, \tilde{y})\) be an approximate eigenpair (either a Ritz or refined Ritz pair) of \( Q_\sigma(\rho)x = 0 \) and \( \tilde{r} = Q_\sigma(\tilde{\rho})y \). Then \((\frac{1}{\rho} + \sigma, y)\) is the corresponding approximate eigenpair of \( Q(\lambda):x = (\lambda^2 M + \lambda C + K)x = 0 \). Define \( \tilde{r} = Q(\frac{1}{\rho} + \sigma)y \). Then we obtain

\[
\frac{\tilde{r}}{\tilde{\rho}^2} = (M_\sigma + C_\sigma/\tilde{\rho} + K_\sigma/\tilde{\rho}^2)y = (\frac{1}{\rho} + \sigma)^2 M + (\frac{1}{\rho} + \sigma)C + K)y = Q(\frac{1}{\rho} + \sigma)y = \tilde{r},
\]

from which it is direct to get the desired \( \|\tilde{r}\| \) from \( \|\tilde{r}\| \) without computing \( \tilde{r} \) explicitly.

We make a final note on Algorithm 5. In previous discussions and analysis, we have supposed \( p = k - m \) previously. This is not mandatory. In order to compute \( m \) desired eigenpairs of (1), the only restriction to \( p \) is that \( p \leq k - m \). So the choice of \( p \) is flexible and takes the form \( p = k - (m + l) \) with \( l \) a very small nonnegative integer, as done in [23] and [12, 13, 15, 16], where \( l = 3 \) is often used. We remark that different \( p \) may have considerable effects on the overall performance of the algorithms, but its choice can only be empirical.

### 4.3. Cure of deflations in implicit restarts

Theorem 2 requires that no deflation occurs in implicit restarts. If deflations occur at steps \( m_1, m_2, \ldots, m_j \leq k \), then the corresponding \( j \) columns \( q_{m_j} \) of \( Q_k \) are zeros. Denote by \( \hat{Q}_k \) and \( \hat{V}_k \) the matrices by deleting the zero columns of \( Q_k \) and rows \( m_1, m_2, \ldots, m_j \) of \( V_k \), respectively. Then we have \( Q_k^+ = \hat{Q}_k \hat{V}_k = \hat{Q}_k \tilde{V}_k \), from which and (9) we get

\[
\begin{bmatrix}
A & B \\
I & 0
\end{bmatrix}
\begin{bmatrix}
\hat{Q}_k \hat{V}_k \\
\hat{P}_k \hat{V}_k
\end{bmatrix} =
\begin{bmatrix}
\hat{Q}_k \tilde{V}_k \\
\hat{P}_k \tilde{V}_k
\end{bmatrix} T_k^+ + t_{k+1} \begin{bmatrix}
q_{k+1} \\
p_{k+1}
\end{bmatrix} e_{k}^* V_k,
\]

where \( T_k^+ = \tilde{V}_k^* T_k V_k \). We see that, although \( \hat{Q}_k \) is still column orthonormal, \( Q_k^+ = \hat{Q}_k \tilde{V}_k \) is not as \( \tilde{V}_k \) is not orthogonal any longer when some rows are deleted from the orthogonal matrix \( V_k \). As a result, \( Q_m^+ = Q_k V_k(:, 1:m) \) is not column orthonormal,
and (17) is not an $m$-step GSOAR decomposition any longer. This means that implicit restarting fails to work whenever deflation occurs.

In what follows we present an effective approach to cure deflation so as to recover a standard GSOAR decomposition, making implicit restarting always applicable unconditionally.

Note that $\hat{V}_k$ is a $(k-j) \times k$ of rank $k-j$. Without loss of generality, we assume that the first $k-j$ columns of $\hat{V}_k$ are linearly independent, i.e., the matrix $\hat{V}_{k1}$ consisting the first $k-j$ columns of $\hat{V}_k$ is nonsingular. Write $\hat{V}_k = [\hat{V}_{k1}, \hat{V}_{k2}]$. We compute the QR decomposition of $\hat{V}_{k1}$ using the Matlab built-in function qr.m and obtain the decomposition of form

\begin{equation}
\hat{V}_k = U_k R_k = [U_{k-j}, 0] \begin{bmatrix} R_{k-j} & R_{12} \\ 0 & I \end{bmatrix},
\end{equation}

where $\hat{V}_{k1} = U_{k-j}R_{k-j}$ is the QR decomposition of $\hat{V}_{k1}$ and $R_{12} = U_{k-j}^* \hat{V}_{k2}$, and $I$ is the identity matrix of order $j$, so that $R_k$ is nonsingular and upper triangular.

Noting that $U_k = \hat{V}_k R_k^{-1}$ and right multiplying (23) by $R_k^{-1}$, we get

\begin{equation}
\begin{bmatrix} A & B \\ I & 0 \end{bmatrix} \begin{bmatrix} \hat{Q}_k U_k \\ P_k V_k R_k^{-1} \end{bmatrix} = \begin{bmatrix} \hat{Q}_k U_k \\ P_k V_k R_k^{-1} \end{bmatrix} R_k T_k^+ R_k^{-1} + t_{k+1k} \begin{bmatrix} q_{k+1} \\ p_{k+1} \end{bmatrix} e_m^* V_k R_k^{-1}.
\end{equation}

Since $R_k^{-1}$ is upper triangular, $R_k T_k^+ R_k^{-1}$ is Hessenberg. Note that $V_k$ has only $p = k - m$ nonzero subdiagonals. Then the first possible nonzero entry $\tilde{\beta}$ of $e_m^* V_k$ is in position $m$ and

\[ t_{k+1k} e_m^* V_k R_k^{-1} = (0, \ldots, 0, \tilde{\beta} b^T) \]

with $\tilde{\beta} = t_{k+1k} V_k(k, m)/e_m^* R_k e_m$. Equating the first $m$ columns on two sides of (25), we obtain

\begin{equation}
\begin{bmatrix} A & B \\ I & 0 \end{bmatrix} \begin{bmatrix} \hat{Q}_m^+ \\ \hat{P}_m^+ \end{bmatrix} = \begin{bmatrix} \tilde{Q}_m^+ \\ \tilde{P}_m^+ \end{bmatrix} T_m^+ + \beta_m^+ \begin{bmatrix} q_{m+1} \\ p_{m+1} \end{bmatrix} e_m^*,
\end{equation}

where $\hat{Q}_m^+ = \hat{Q}_k U_k(:, 1:m)$, $\hat{P}_m^+ = P_k V_k(:, 1:m) R_m^{-1}$ with $R_m$ the $m \times m$ leading principal matrix of $R_k$, $T_m^+$ the $m \times m$ leading principal matrix of $R_k T_k^+ R_k^{-1}$, and

\begin{equation}
\begin{bmatrix} q_{m+1} \\ p_{m+1} \end{bmatrix} = \frac{1}{\beta_m} \begin{bmatrix} f_m^+ \\ g_m^+ \end{bmatrix} = \begin{bmatrix} \hat{Q}_k U_k \\ P_k V_k R_k^{-1} \end{bmatrix} e_m^+ + \tilde{\beta} \begin{bmatrix} q_{k+1} \\ p_{k+1} \end{bmatrix},
\end{equation}

\begin{equation}
\beta_m^+ = \|\hat{f}_{m+1}^+ \hat{Q}_k U_k e_m^+ + \tilde{\beta} q_{k+1}\|.
\end{equation}

(24) indicates that the column orthonormality of $U_k(:, 1:m)$ is guaranteed whenever $m \leq k-j$, i.e., $j \leq k-m$. This means that $\hat{Q}_m^+ = \hat{Q}_k U_k(:, 1:m)$ is column
orthonormal, provided that the number $j$ of deflations during the last cycle of GSOAR procedure does not exceed $k - m$. If $m > k - j$, the first $k - j$ columns of $Q^+_m$ are orthonormal and the last $m - (k - j)$ columns of $U_k$ are zero, so that the last $m - (k - j)$ columns of $Q^+_m$ are zero. As a result, there are $m - (k - j)$ deflations in (26). For either $m \leq k - j$ or $m > k - j$, it is trivial to justify that $(Q^+_m)q^+_m = 0$. Therefore, by curing deflations in implicit restarts, we have obtained a truly $m$-step GSOAR decomposition (26).

5. Numerical Experiments

In this section we report numerical examples to illustrate the practicability of IGSOAR and IRGSOAR and the superiority of IRGSOAR to IGSOAR. Meanwhile, we also compare them with the corresponding counterparts IRSGA and IRRSGA proposed in [10] for some test problems. In addition, we compare IGSOAR and IRGSOAR with the Matlab function \texttt{eigs}, the implicitly restarted Arnoldi method with exact shifts used, which is directly applied to the linearization problem (3). All the experiments were run on Intel(R)Core(TM)i5-3470s CPU 2.9GHz, RAM 4G using Matlab R2012b with $\epsilon_{\text{mach}} = 2.22 \times 10^{-16}$ under the Windows 7 system.

We list CPU timings (in second) of the three main parts abbreviated as ‘SOAR’, ‘SMALL’ and ‘IMRE’, where ‘SOAR’ denotes the CPU time of the first cycle of GSOAR procedure plus standard extensions of the GSOAR decomposition from step $m + 1$ to step $k$ for all the other cycles, ‘SMALL’ is the CPU time of forming the projected QEP, solving them and computing residuals of approximate eigenpairs, and ‘IMRE’ is the CPU time of performing all implicit QR iterations and generating the $m$-step GSOAR decompositions for all cycles. In addition, we use ‘restarts’ and ‘CPU time’ to denote the number of restarts and the total CPU time of IGSOAR, IRGSOAR and \texttt{eigs}, respectively.

For each example, we used the same starting vector generated randomly in a uniform distribution for IGSOAR and IRGSOAR. We transformed the projected QEP (12) to the generalized eigenvalue problem (19) and solved it by the QZ algorithm, i.e., the Matlab built-in function \texttt{eig.m}. We recovered an eigenvector $g$ of QEP (12) from either the first $k$ components or the last $k$ components of $[\theta g^T, q^T]^T$. From the backward error analysis [7], it is preferable to take the first $k$ ones if $|\theta| \geq 1$ and the last $k$ ones if $|\theta| < 1$. We adopted this choice.

For \texttt{eigs}, we used the same $k$ as that in IGSOAR and IRGSOAR to compute the same $m$ eigenpairs for each example. The CPU time of \texttt{eigs} did not include the time of computing the LU decomposition of $M$, which is used when acting a matrix-vector product in \texttt{eigs} at each step. The starting vector of \texttt{eigs} was obtained by normalizing $\begin{bmatrix} q_1 \\ p_1 \end{bmatrix}$,
where $q_1$ and $p_1$ were the vectors in Algorithm 3. The number of shifts was the default value, i.e., $p = k - (m + 3)$. We also used $tol$ to denote the stopping criterion used in eigs for (3). Let $(\theta, y)$ be a converged eigenpair computed by eigs, we set $y_1$ to be the vector consisting of the first $n$ components of $y$, and $y_2$ the vector consisting of the last $n$ components of $y$. We then computed the relative residual norms (20) of $(\theta, y_1)$ and $(\theta, y_2)$ and took the smaller one as the residual norm of eigs for QEP (1). $\text{'Res}_{\text{min}}$ ' and $\text{'Res}_{\text{max}}$ ' recorded the minimum and maximum relative residual norms (20) obtained in this way for all the converged eigenpairs for (3). The maximum number of restarts is limited to 50.

Example 1. We consider the damped vibration mode of an acoustic fluid confined in a cavity with absorbing walls capable of dissipating acoustic energy [9]. We take the same geometrical data as in [9]. The QEP is

$$\lambda^2 M_u u + (\alpha + \lambda \beta) A u + K u = 0,$$

where $\alpha = 5 \times 10^4 N/m^3$, $\beta = 200 N s/m^3$, and the order $n = 46548$.

By taking $tol = 10^{-14}$ and two sets of parameters $k = 30, p = 7$ and $k = 30, p = 5$, we used IRGSOAR and IGSOAR to compute the twenty eigenvalues nearest to the complex target $\sigma = 25 + 18\pi i$ and the corresponding eigenvectors of the above QEP. Table 1 reports the results obtained, and Figure 1 describes the convergence processes of two algorithms, depicting the maxima of relative residual norms of $m$ approximate eigenpairs versus restarts.

We see from Table 1 and Figure 1 that two algorithms were efficient. However, as far as both restarts and CPU timings are concerned, IRGSOAR was twice as fast as IGSOAR for $k = 30$ and $p = 5$, and the former was also considerably faster than the latter for $k = 30$ and $p = 7$. Furthermore, we observe from the figure that the residual norm of IRGSOAR was smaller than that of IGSOAR substantially at each cycle, indicating that the refined Ritz vectors can be considerably more accurate than the Ritz vectors. We find that for the same $k$, the value of $p$ has an effect on the overall performance of IGSOAR and IRGSOAR. For this example, we took two $p$ smaller than $k - m = 10$. It is seen that the effect is marginal for IRGSOAR, while it is relatively essential for IGSOAR. In addition, we remark that the most consuming cost was paid to the SOAR procedure, but the explicit computation and solutions of all small QEP also occupied quite portion of the CPU time cost. The CPU time 'IMRE' of implicit restarting consumed least but could not be negligible.

For this example, by taking $tol = 10^{-6}$, we found that eigs was much more costly than IGSOAR and IRGSOAR to converge and the approximate eigenpairs were as accurate as those obtained by the latter two algorithms, while, for $tol = 10^{-8}$, it failed to converge after 50 restarts. We should point out that our codes are programmed in the Matlab language and may not be optimized while eigs is programmed in C language and optimized. This means that for the same $k$ each restart of eigs should be
more time consuming than that of IGSOAR and IRGSOAR since \texttt{eigs} is much more expensive than IGSOAR and IRSOAR in the orthogonalization of Arnoldi vectors. As a result, in all the experiments the number of restarts is more reasonable to compare the computational efficiency of these three algorithms. It is worthwhile to mention that for this example a relatively big $tol = 10^{-6}$ for \texttt{eigs} delivered very accurate eigenpairs of QEP (1) and a smaller $tol$ is unnecessary.

Table 1: Example 1, $tol=10^{-14}$

| Algorithm | $k$ | $p$ | restarts | CPU time | SOAR | SMALL | IMRE |
|-----------|----|----|---------|---------|------|-------|------|
| IRGSOAR   | 30 | 7  | 3       | 12.55   | 8.56 | 2.58  | 1.34 |
| IGSOAR    | 30 | 5  | 5       | 16.86   | 10.99| 3.12  | 2.68 |
| IRGSOAR   | 30 | 7  | 3       | 11.96   | 7.93 | 2.58  | 1.31 |
| IGSOAR    | 30 | 7  | 5       | 21.01   | 12.06| 4.57  | 4.27 |

The results obtained by \texttt{eigs}

| $tol$ | $k$ | CPU | restarts | $Res_{min}$ | $Res_{max}$ |
|-------|----|-----|----------|--------------|--------------|
| $10^{-6}$ | 30 | 27.89 | 15      | $0.42 \times 10^{-15}$ | $0.93 \times 10^{-14}$ |
| $10^{-8}$ | 30 | –    | 50      | –            | –            |

Fig. 1. Example 1. Residuals versus restarts. Left: $k=30$, $p=5$; right: $k=30$, $p=7$.

Example 2. This problem and arises in a model of the concrete structure supporting a machine assembly [4, 5] and has the form $Q(\lambda)x = (\lambda^2 M + \lambda C + (1 + i\mu)K)x = 0$. The matrices are of order 2472, where $M$ is real diagonal, $C$, the viscous damping matrix, is pure imaginary and diagonal, $K$ is complex symmetric, and the factor $1 + i\mu$ adds uniform hysteretic damping. We use the command \texttt{nlevp(‘concrete’, 0.04)} in [4] to generate the complex symmetric coefficient matrices. Thus problem was tested in [10].

We ran IRGSOAR and IGSOAR to compute the ten eigenvalues nearest to the origin by taking $tol = 10^{-14}$ and the same $k = 20$, two $p = 7$ and 5. Table 2 and Figure 2 reported the results, from which it can be seen that two algorithms worked very
well and IRGSOAR was a little more efficient than IGSOAR in terms of both restarts and CPU timings. We remark that, for this problem, the corresponding algorithms IRSGA and IRRSGA in [10] both used four restarts to achieve the convergence for the same $k = 20$ and $tol$. Note that they use the F-norm in the denominator of (20), which means that for the same $tol$ our convergence tolerance is smaller. Therefore, for $p = 7$, IGSOAR was (at least) as efficient as IRSGA, and IRGSOAR was faster than IRRSGA. For $p = 5$, IGSOAR used five restarts for a smaller stopping tolerance than that used by IRSGA, and IRGSOAR used four restarts. This demonstrates that, for this problem, IGSOAR and IRGSOAR were as efficient as IRSGA and IRRSGA, respectively. It is clear that two different $p$ affected the overall efficiency of each algorithm only marginally. Finally, we observe that, unlike Example 1, the main cost of each algorithm was paid to the GSOAR procedure and overwhelmed "SMALL" and "IMRE".

Table 2: Example 2, $tol = 10^{-14}$

| Algorithm | $k$ | $p$ | restarts | CPU time | SOAR | SMALL | IMRE |
|-----------|-----|-----|----------|----------|------|-------|------|
| IRGSOAR   | 20  | 7   | 3        | 0.78     | 0.62 | 0.10  | 0.03 |
| IGSOAR    | 20  | 7   | 4        | 1.03     | 0.91 | 0.07  | 0.04 |
| IRGSOAR   | 20  | 5   | 4        | 0.90     | 0.72 | 0.12  | 0.05 |
| IGSOAR    | 20  | 5   | 5        | 0.93     | 0.77 | 0.09  | 0.07 |

The results obtained by eigs

| $tol$ | $k$ | CPU time | restarts | $Res_{min}$ | $Res_{max}$ |
|-------|-----|----------|----------|-------------|-------------|
| $10^{-5}$ | 20  | 1.03 | 7        | $0.19 \times 10^{-14}$ | $0.13 \times 10^{-13}$ |
| $10^{-10}$ | 20  | 1.25 | 9        | $0.19 \times 10^{-14}$ | $0.12 \times 10^{-13}$ |
| $10^{-14}$ | 20  | 1.70 | 12       | $0.13 \times 10^{-14}$ | $0.52 \times 10^{-17}$ |

Fig. 2. Example 2. Residuals versus restarts; Left: $k = 20, p = 7$; right: $k = 20, p = 5$.

We also report the results obtained by eigs for three $tol$ and list them in Table 2. It is seen from Table 2 that both IGSOAR and IRGSOAR performed much better than
eigs, and they used much less CPU time and fewer restarts to compute the desired eigenpairs with much higher accuracy. As $\text{Res}_{\text{min}}$ and $\text{Res}_{\text{max}}$ indicated, the accuracy of the converged eigenpairs obtained by eigs with three greatly varying tol essentially had no difference as the approximate eigenpairs of QEP (1), and their relative residual norms were already at the level of $\epsilon_{\text{mach}}$ for $\text{tol} = 10^{-8}$.

Example 3. This example is from [4] and tested in [10] (cf. Example 6.3). We tested IRGSOAR and IGSOAR for the following cases (a) and (b) by taking $\text{tol} = 10^{-14}$.

Case (a): Acoustic 1D. This example arises from the finite element discretization of the time harmonic wave equation $-\Delta p - (2\pi f/c)^2 p = 0$. Here, $p$ denotes the pressure, $f$ is the frequency, $c$ is the speed of sound in the medium, and $\xi$ is the (possibly complex) impedance. On the domain $[0, 1]$ with $c = 1$, the $n \times n$ matrices $M$, $D$, and $K$ are defined by

$$M = -4\pi^2 \frac{1}{n} (I - e_n e_n^T), \quad D = 2\pi i \frac{1}{\xi} e_n e_n^T, \quad K = n (\text{tridiag}(-1, 2, -1) - e_n e_n^T).$$

### Table 3: Example 3(a), tol=$10^{-14}$

| Algorithm | $k$ | $p$ | restarts | CPU time | SOAR | SMALL | IMRE |
|-----------|-----|-----|----------|----------|------|-------|------|
| IRGSOAR   | 12  | 5   | 2        | 0.12     | 0.02 | 0.06  | 0.02 |
| IRGSOAR   | 12  | 5   | 3        | 0.12     | 0.02 | 0.05  | 0.03 |
| IGSOAR    | 12  | 3   | 2        | 0.10     | 0.02 | 0.06  | 0.01 |
| IGSOAR    | 12  | 3   | 3        | 0.10     | 0.02 | 0.05  | 0.02 |

The results obtained by eigs

| tol       | $k$ | CPU time | restarts | $\text{Res}_{\text{min}}$ | $\text{Res}_{\text{max}}$ |
|-----------|-----|----------|----------|---------------------------|---------------------------|
| $10^{-8}$  | 12  | 0.27     | 10       | $0.32 \times 10^{-18}$    | $0.97 \times 10^{-13}$    |
| $10^{-10}$ | 12  | 0.72     | 22       | $0.40 \times 10^{-18}$    | $0.20 \times 10^{-13}$    |
| $10^{-14}$ | 12  | 0.75     | 23       | $0.81 \times 10^{-18}$    | $0.83 \times 10^{-18}$    |

Fig. 3. Example 3(a). Left: $k = 12$, $p = 5$; right: $k = 12$, $p = 3$. 
We use `nlevp(acoustic_wave_1d,5000,1)` to generate matrices $M$, $D$ and $K$ with size $n = 5000$.

Just as in [10], we computed the six eigenvalues nearest to the origin with $k = 12$, $p = 5$ and 3. Table 3 reports the results, and Figure 3 depicts the convergence processes of two algorithms. From the figure we see that, for the same $k$ and two $p$, IRGSOAR and IGSOAR used two and three cycles, respectively. As indicated in [10], for the same $k = 12$ and a little larger convergence tolerance, the corresponding implicitly restarted algorithms IRSGA and IRRSGA both used three cycles. So IRGSOAR was a little better than IRRSGA. Regarding CPU time, since $M$, $D$ and $K$ are very sparse, the CPU timings of the GSOAR procedure and implicit restarting are comparable, and less than ‘SMALL’.

As we have seen, the eigenpairs obtained by `eigs` had similar accuracy to those obtained by IGSOAR and IRGSOAR with three greatly varying $tol$, and all of them converged to the level of machine precision. However, Table 3 clearly shows that IGSOAR and IRGSOAR were much more efficient than `eigs`.

**Table 4: Example 3(b), $tol=10^{-14}$**

| Algorithm | $k$ | $p$ | restarts | CPU time | SOAR | SMALL | IMRE |
|-----------|-----|-----|-----------|----------|------|-------|------|
| IRGSOAR   | 12  | 5   | 7         | 1.21     | 0.59 | 0.39  | 0.19 |
| IGSOAR    | 12  | 5   | 11        | 1.57     | 0.85 | 0.36  | 0.33 |
| IRGSOAR   | 12  | 3   | 10        | 1.18     | 0.45 | 0.45  | 0.24 |
| IGSOAR    | 12  | 3   | 10        | 0.94     | 0.46 | 0.25  | 0.22 |

The results obtained by `eigs`:

| $tol$  | $k$ | $CPU$ time | restarts | $Res_{min}$     | $Res_{max}$    |
|--------|-----|------------|----------|-----------------|----------------|
| $10^{-8}$ | 12 | 0.84       | 8        | $0.38 \times 10^{-17}$ | $0.48 \times 10^{-13}$ |
| $10^{-10}$ | 12 | 0.83       | 8        | $0.40 \times 10^{-17}$ | $0.14 \times 10^{-13}$ |
| $10^{-14}$ | 12 | 0.99       | 11       | $0.32 \times 10^{-17}$ | $0.53 \times 10^{-16}$ |

![Fig. 4. Example 3(b). Left: $k = 12$, $p = 5$; right: $k = 12$, $p = 3$.](image)

**Case (b):** Acoustic 2D. This example is a two-dimensional acoustic wave equation on $[0, 1] \times [0, 1]$. The coefficient matrices $M$, $D$ and $K$ are given by
\[ M = -4\pi^2 h^2 I_{q-1} \otimes \left( I_q - \frac{1}{2} e_q e_q^T \right), \quad D = 2\pi i \frac{h}{\xi} I_{q-1} \otimes (e_q e_q^T), \]
\[ K = I_{q-1} \otimes D_q + T_{q-1} \otimes \left( -I_q + \frac{1}{2} e_q e_q^T \right). \]

where \( h \) denotes the mesh size, \( q = 1/h \), \( \otimes \) denotes the Kronecker product, \( \xi \) is the (possibly complex) impedance, \( D_q = \text{tridiag}(-1, 4, -1) - 2e_q e_q^T \), and \( T_{q-1} = \text{tridiag}(1, 0, 1) \). We use \texttt{nlevp(‘acoustic\_wave\_2d’,90,0.1*1i)} to get the real symmetric matrices \((M, D, K)\). The matrix size is given by \( n = 8010 \).

As in [10], we computed the six eigenvalues nearest to the origin with \( k = 12, p = 5 \) and 3. Table 4 and Figure 4 give the results and convergence processes of two algorithms, respectively. It is seen that IRGSOAR and IGSOAR used seven and eleven restarts for \( p = 5 \), respectively, and both of them used ten cycles for \( p = 3 \). Therefore, two algorithms were efficient, and IRGSOAR could be more efficient than IGSOAR. We find that both IGSOAR and IRGSOAR were more efficient than IRSGA and IR-RSGA [10], where the latter ones used eleven and twelve cycles, respectively.

For this problem, unlike the previous examples, \texttt{eigs} performed very well and was comparable to IGSOAR and IRGSOAR in terms of the accuracy and the computational efficiency.

**Example 4.** This QEP arises in an \( n \)-degree-of-freedom damped mass-spring system [25]. By taking \( m_i = 1 \) and letting all the springs (respectively, dampers) have the same constant \( \kappa \) (respectively, \( \tau \)) except \( \kappa_1 = \kappa_n = 2\kappa \) and \( \tau_1 = \tau_n = 2\tau \), the resulting matrices are
\[ M = I, \quad C = \tau \cdot \text{tridiag}(-1, 3, -1), \quad K = \kappa \cdot \text{tridiag}(-1, 3, -1), \]
which are very sparse. We took \( n = 5000, \kappa = 5 \) and \( \tau = 10 \) and were interested in the six eigenvalues nearest to the complex target \( \sigma = -13 + 0.4i \) and the corresponding eigenvectors.

For \( tol = 10^{-10} \), we tested IRGSOAR and IGSOAR for \( k = 40, p = 23 \) and 28. Table 5 lists the results, and Figure 5 depicts the convergence processes for two sets of parameters \( k \) and \( p \).

It can be found from Table 5 and Figure 5 that two algorithms worked quite well. Compared with Examples 1-3, much more restarts were needed now; for the given \( k \), two different \( p \) did not make much difference on restarts and CPU timings of two algorithms. Furthermore, IRGSOAR and IGSOAR are similarly efficient, and the former used a little fewer restarts but more CPU time than IGSOAR. Since the matrices in this QEP are very sparse, it appears that performing the SOAR procedure in each algorithm was not dominant, and instead it was considerably less costly than the explicit computation and solutions of all small QEP and implicit restarting, as indicated by Table 5.
Table 5: Example 4, \( tol=10^{-10} \)

| Algorithm  | \( k \) | \( p \) | restarts | CPU time | SOAR  | SMALL | IMRE |
|------------|--------|--------|----------|----------|-------|-------|------|
| IRGSOAR    | 40     | 23     | 41       | 11.01    | 1.78  | 4.74  | 4.39 |
| IGSOAR     | 40     | 23     | 44       | 9.02     | 1.88  | 2.33  | 4.73 |
| IRGSOAR    | 40     | 28     | 39       | 9.96     | 1.79  | 4.07  | 4.01 |
| IGSOAR     | 40     | 28     | 47       | 9.50     | 2.18  | 2.47  | 4.78 |

The results obtained by \( \texttt{eigs} \)

| \( tol \) | \( k \) | CPU time | restarts | \( Res_{\min} \) | \( Res_{\max} \) |
|-----------|--------|----------|----------|----------------|----------------|
| \( 10^{-8} \) | 40     | 6.05     | 31       | \( 0.37 \times 10^{-5} \) | \( 0.37 \times 10^{-5} \) |
| \( 10^{-10} \) | 40     | 9.36     | 47       | \( 0.37 \times 10^{-5} \) | \( 0.37 \times 10^{-5} \) |

Fig. 5. Example 4. Residuals versus restarts. Left: \( k = 40, p = 23 \); right: \( k = 40, p = 28 \).

For this example, unlike all the previous examples, for given two \( tol \) similar to that used by IGSOAR and IRGSOAR, \( \texttt{eigs} \) used comparable restarts and the CPU time to IGSOAR and IRGSOAR, but it computed the desired eigenpairs with much poorer accuracy, so, as a whole, it is considerably inferior to IGSOAR and IRGSOAR. An important observation is that improving the accuracy of approximate eigenpairs of (3) may be helpless to improve their accuracy as the approximate eigenpairs of (1). A comparison of this example and Example 1 reveals a remarkable difference: \( \texttt{eigs} \) with big \( tol \) computed the desired eigenpairs with the accuracy at the level of \( \epsilon_{\text{mach}} \) for Example 1, while it with smaller \( tol \) got the desired eigenpairs with much poorer accuracy. So it is uncertain for us to choose a suitable \( tol \) for \( \texttt{eigs} \) to compute the desired eigenpairs with a prescribed accuracy in the sense of the stopping criterion (20) for QEP (1).

Example 5. This problem comes from [4]. It is a nonlinear eigenvalue problem modeling a radio-frequency gun cavity that is of the form

\[
T(\lambda)x = [K - \lambda M + i(\lambda - \sigma_1^2)^{1/2}W_1 + i(\lambda - \sigma_2^2)^{1/2}W_2]x = 0,
\]
where $M$, $K$, $W_1$, $W_2$ are real symmetric matrices of size $9956 \times 9956$. From these matrices, we constructed a QEP of the form

$$(\lambda^2W_2 + \lambda M + K)x = 0,$$

which is purely for our test purpose. We used IRGSOAR and IGSOAR to compute the six eigenvalues nearest to $\sigma = 0.5 + 0.5i$ and the associated eigenvectors. Table 6 and Figure 6 reported the results.

**Table 6: Example 5, tol=10^{-10}**

| Algorithm | $k$ | $p$ | restarts | CPU time | SOAR | SMALL | IMRE |
|----------|----|----|---------|---------|------|-------|------|
| IRGSOAR  | 20 | 5  | 1       | 1.73    | 1.59 | 0.11  | 0.00 |
| IGSOAR   | 20 | 5  | 9       | 6.10    | 4.93 | 0.58  | 0.54 |
| IRGSOAR  | 20 | 11 | 1      | 1.73    | 1.61 | 0.11  | 0.00 |
| IGSOAR   | 20 | 11 | 3      | 3.77    | 3.42 | 0.19  | 0.13 |

The results obtained by eigs

| $tol$ | $k$ | CPU time | restarts | $Res_{\text{min}}$ | $Res_{\text{max}}$ |
|-------|----|----------|---------|-------------------|-------------------|
| $10^{-6}$ | 20 | 2.73 | 4 | $0.37 \times 10^{-13}$ | $0.45 \times 10^{-8}$ |
| $10^{-8}$ | 20 | 59.31 | 50 | $0.25 \times 10^{-7}$ | $0.21 \times 10^{-3}$ |

*Fig. 6. Example 5. Residuals versus restarts. Left: $k=20$, $p=5$; right: $k=20$, $p=11$.*

For this example, two algorithms worked well. However, IRGSOAR exhibited the very considerable superiority to IGSOAR. We find the desired eigenpairs without restarting the algorithm for given two sets of parameters $k$ and $p$ while IGSOAR used nine and three cycles, respectively. In terms of CPU timings, IRGSOAR was also a few times faster than IGSOAR. Furthermore, for this example, the CPU time of the SOAR procedure dominated the CPU time cost of each algorithm. On contrary to Example 4, for the given $k$, the smaller $p = 5$ made IGSOAR use considerably more restarts and CPU time, meaning that the choice of $p$ may have considerable effects on the overall performance of IGSOAR. However, this example and Examples 1-4 illustrate that the
effects of $p$ must be problem dependent, and it is impossible to design a definite and general effective way to select it.

In contrast, \texttt{eigs} behaved not good for this example, and it used much more restarts to achieve the convergence for the not much smaller $tol = 10^{-8}$ than $10^{-6}$. However, as approximate eigenpairs of QEP (1), the converged eigenpairs with $tol = 10^{-6}$ were substantially more accurate than those with $tol = 10^{-8}$. This is really bad because it shows that, on the contrary to our common acceptance, that considerably more accurate eigenpairs for the linearization problem (3) are not necessarily more accurate too for QEP (1). This, together with Example 1 and Example 4, demonstrates that solving the linearization problem (3) directly has serious uncertainty, as far as the accuracy is concerned.

6. CONCLUSION

We have considered generalized second-order Arnoldi method and its refined version for solving the large QEP. The methods are structure-preserving and applied to the QEP directly after an orthonormal basis of the generalized second-order Krylov subspace is generated by the GSOAR procedure. To be practical, we have developed implicitly restarted algorithms with certain exact and refined shifts proposed for two methods, respectively. We have presented an efficient and reliable algorithm for computing the shift candidates. Unlike Arnoldi type algorithms for the linear eigenvalue problem, where the number of shift candidates are just that of shifts, for the QEP the shift candidates are more than the shifts. we have discussed in detail how to seek and determine reasonable shifts for each method. Also, deflation may occur in the algorithm for the QEP, for which implicit restarting is not applicable. To overcome this deficiency, we have proposed an effective approach to cure deflation in implicit restarts, so that implicit restarting can be used to the GSOAR procedure unconditionally. We have tested our algorithms on a number of real-world problems. Numerical experiments have demonstrated that two algorithms work well and the refined algorithm can outperform the standard counterpart considerably. They also show that our algorithms generally perform much better than \texttt{eigs} in terms of the accuracy or the computational efficiency.

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