Abstract. The article provides an explicit algebraic expression for the generating function of walks on graphs. Its proof is based on the scattering theory for the differential Laplace operator on non-compact graphs.

1. Introduction

The concept of a generating function is known to be a very important tool in combinatorics, probability, and number theory. Associated methods reduce the solution of combinatorial or probabilistic problems to the study of particular properties of the generating function which can be performed by methods of function theory and analysis. For an introduction to this method the reader may consult the books [6], [8], [32]. A number of solved and still unsolved combinatorial problems, where the generating function plays a central role can be found in the article [25]. In the probabilistic context we mention the solution of the problem whether a simple random walk on $\mathbb{Z}^d$ is recurrent or transitive by an analysis of the generating function (see, e.g., [9]). Some further examples will be discussed below in Sections 2 and 7.

The present work is devoted to the determination of the generating function for walks on graphs (both in combinatorial and probabilistic contexts). Walks on graphs are considered, in particular, in [2], [5], [18], [26, Section 4.7], [33]. In a custom setting random walks on graphs are defined as Markov chains on the vertices of the graph. The transition probability from one vertex to another is assumed to be non-zero if and only if these vertices are adjacent. For a survey of the theory of random walks on graphs, see [17].

We consider a slightly different but closely related model of random walks on graphs, where the states are chosen to be the edges of the graph. Transitions between different states are determined by stochastic (that is, Markov) matrices $M(v)$ prescribed at every vertex of the graph. The graphs are assumed to be non-compact, that is, besides a finite number of edges (or “internal lines” $i \in \mathcal{I}$) to have a non-empty set $\mathcal{E}$ of “external lines” which serve as entries or exits for random walks. More precisely the model will be described in the following section. A relation between this model of random walks and random walks on vertices is explained in the Appendix below.

Consider an arbitrary positive weight on the graph (that is, a map assigning to any edge $i$ of the graph a positive number $a_i$). We will call $a = \{a_i\}_{i \in \mathcal{I}} \in (\mathbb{R}_+)^{\mathcal{I}}$ a penalty vector. With $\beta$ being a complex parameter we define a generating function $T_{e,e'}(\beta)$ of
walks from an external line \(e' \in E\) to an external line \(e \in E\) as

\[
T_{e,e'}(\beta) = \sum [M(v_N)]_{e,i_N} e^{-\beta a_{i_N}} [M(v_{N-1})]_{i_N,i_{N-1}} \cdots [M(v_1)]_{i_2,i_1} e^{-\beta a_{i_1}} [M(v_0)]_{i_1,e'},
\]

where the sum is taken over all walks \(\{e',i_1,i_2,\ldots,i_{N-1},i_N,e\}\) from \(e'\) to \(e\), the set \(v_0,v_1,\ldots,v_N\) is the ordered list of vertices (with possible repetitions) visited by the walk, \(i_k\) are the corresponding internal lines traversed during the walk (again with possible repetitions). In the context of the generating functions the weight \(\exp\{-\beta a_i\}\) can be viewed as a penalty factor for traversing the edge \(i\) during a walk.

More generally, one can also consider penalty vectors depending on the direction in which a given edge is traversed by the walk. The corresponding generating function will be discussed in Section 6 below (see Theorem 6.7).

The main result of the present work (see Theorem 6.2 below) provides an explicit algebraic expression for the generating function of walks on graphs. Its proof is based on the scattering theory for the differential Laplace operator on non-compact graphs and the corresponding methods developed by the authors in \([12],[13],[14],[15],[16]\). In the context of differential operators the weights \(a_i\) will be interpreted as the metric lengths of the edges \(i\).

The generating function is determined by analytic continuation of the scattering matrix to complex values of the spectral parameter. This result is very reminiscent of a similar result in relativistic quantum field theory in the context of vacuum expectations of products of quantum fields. The analytic continuation of the Wightman distributions \([31]\) to the Euclidean points (the so called Wick rotation \([30]\)) results in the Schwinger functions \([24]\). Conversely, by a result \([20],[21]\) of K. Osterwalder and one of the authors (R. S.) the Schwinger functions give rise to Wightman distributions. In the bosonic case Symanzik and Nelson have shown that the Schwinger functions describe a stochastic theory (see \([7],[19],[28]\), and references quoted there).

We expect that the model of walks on graphs considered in the present article may be of interest in the context of optimization of traffic flows and in telecommunication networks, where the transition matrices \(M(v)\) determine a proportion of the traffic or signals to be transmitted in a given direction.

The article is organized as follows. In Section 2 we will give definitions of walks on graphs and of associated generating functions. Also we present several examples relating the generating function to combinatorics. In Section 3 we will revisit the scattering theory of differential Laplace operators on graphs. Section 4 is devoted to the proof of the combinatorial Fourier expansion formula \((4.18)\). Theorem 4.2 proves the absolute convergence of the Fourier series and Theorem 4.10 expresses the Fourier coefficients as sums over the walks on the graph. In Section 5 we will consider the analytic continuation of the scattering matrix with respect to the square root of the energy (that is, the spectral parameter). In Section 6 the generating function will be expressed in terms of the scattering matrix for a Laplace operator with boundary conditions determined by the transition matrices \(M(v)\). In Section 7 we will turn to random walks on graphs. By means of the generating function we will calculate several mean values associated to this probabilistic set-up.

There are several further models which can also be treated by the methods of the present work. In particular, choosing the matrices \(M(v)\) as independent random variables one obtains a model of random walks in random environment. Further, the graph itself can be chosen to be random (see, e.g., \([4]\)). We note that random graphs have been
used to model the spread of epidemics like AIDS, see, e.g., the article \cite{3} and further references quoted there.

2. Walks on Graphs

We consider a finite, connected and non-compact graph $G = (V, I, E, \partial)$, where $V = V(G)$ is a finite set of vertices, $I$ is a finite set of internal lines, $E$ is a finite set of external lines. The elements of the set $I \cup E$ are called edges. The boundary operator $\partial$ assigns to each internal line $i \in I$ an ordered pair $(v_1, v_2)$ of vertices (possibly equal) and to each external line $e \in E$ a single vertex $v$. The vertices $v_1 := \partial^-(i)$ and $v_2 := \partial^+(i)$ are called the initial and terminal vertex of the internal line $i$, respectively. This obviously induces an orientation on each of the internal lines and this will become relevant below.

The vertex $v = \partial(e)$ is the initial vertex of the external line $e$. If $\partial(i) = (v, v)$, then $i$ is called a tadpole. To simplify the discussion, in what follows we will assume that the graph $G$ contains no tadpoles.

Two vertices $v$ and $v'$ are called adjacent if there is an internal line $i \in I$ such that either $(v, v') = \partial(i)$ or $(v', v) = \partial(i)$. A vertex $v$ and the (internal or external) line $j \in I \cup E$ are incident if $v \in \partial(j)$. The degree $\deg(v)$ equals the number of (internal or external) lines incident with the vertex $v$.

We do not require the map $\partial : I \rightarrow V \times V$, $E \rightarrow V$ to be one-to-one. In particular, any two vertices are allowed to be adjacent to more than one internal line and two different external lines may be incident with the same vertex.

Given an arbitrary vector $a = \{a_i\}_{i \in I} \in \mathbb{R}^{|I|}$ with strictly positive components, we will endow the graph with the following metric structure. Any internal line $i \in I$ will be associated with an interval $[0, a_i]$ with $a_i > 0$ such that the initial vertex of $i$ corresponds to $x = 0$ and the terminal one - to $x = a_i$. Any external line $e \in E$ will be associated with a half-line $[0, +\infty)$. The number $a_i$ can be viewed as the length of the internal line $i$.

A nontrivial walk $w$ on the graph $G$ from $e' \in E$ to $e \in E$ is a sequence

$$\{e', i_1, \ldots, i_N, e\}$$

of edges such that

(i) $v_0 := \partial(e') \in \partial(i_1)$, $v_N := \partial(e) \in \partial(i_N)$, and for any $k \in \{1, \ldots, N - 1\}$ there is a vertex $v_k \in V$ such that $v_k \in \partial(i_k)$ and $v_k \in \partial(i_{k+1})$;

(ii) $v_k \neq v_{k+1}$ for all $k \in \{0, \ldots, N - 1\}$.

The number $N$ is the combinatorial length $|w|_{\text{comb}} \in \mathbb{N}$ and the number

$$|w| = \sum_{k=1}^{N} a_{i_k} > 0$$

is the metric length of the walk $w$.

Example 2.1. Let $G = (V, I, E, \partial)$ with $V = \{v_0, v_1\}$, $I = \{i\}$, $E = \{e\}$, $\partial(e) = v_0$, and $\partial(i) = (v_0, v_1)$. Then the sequence $\{e, i, e\}$ is not a walk, whereas $\{e, i, i, e\}$ is a walk from $e$ to $e$.

Proposition 2.2. Given an arbitrary nontrivial walk $w = \{e', i_1, \ldots, i_N, e\}$ there is a unique sequence $\{v_k\}_{k=0}^{N}$ of vertices such that $v_0 = \partial(e') \in \partial(i_1)$, $v_N = \partial(e) \in \partial(i_N)$, $v_k \in \partial(i_k)$, and $v_k \in \partial(i_{k+1})$. 
We emphasize, that at any vertex of the sequence \(\{v_k\}_{k=0}^N\) associated with a nontrivial walk \(w\), the walk is either “reflected” or “transmitted”.

A trivial walk \(w\) on the graph \(G\) from \(e' \in E\) to \(e \in E\) is the tuple \(\{e', e\}\) with \(\partial(e) = \partial(e')\). Both the combinatorial and the metric length of a trivial walk are zero.

Let \(\mathcal{W}_{e,e'} = \mathcal{W}_{e,e'}(G)\), \(e, e' \in E\) be the set of all walks \(w\) on \(G\) from \(e'\) to \(e\). In particular, the set \(\mathcal{W}_{e,e'}\) is infinite for all \(e, e' \in E\) if \(\mathcal{I} \neq \emptyset\) and the graph \(G\) is connected. By reversing a walk \(w\) from \(e'\) to \(e\) into a walk \(w_{rev}\) from \(e\) to \(e'\) we obtain a natural one-to-one correspondence between \(\mathcal{W}_{e,e'}\) and \(\mathcal{W}_{e',e}\). Obviously, \(|w| = |w_{rev}|\) and \(\overline{w}(w) = \overline{w}(w_{rev})\).

Let \(\mathcal{S}(v) \subseteq E \cup \mathcal{I}\) denote the star graph of the vertex \(v \in V\), i.e., the set of the edges adjacent to \(v\). Also, by \(\mathcal{S}_-(v)\) (respectively \(\mathcal{S}_+(v)\)) we denote the set of the edges for which \(v\) is the initial vertex (respectively terminal vertex). Obviously, \(\mathcal{S}_+(v) \cap \mathcal{S}_-(v) = \emptyset\) since \(G\) does not contain tadpoles by assumption.

To every \(v \in V\) we associate an arbitrary \(\deg(v) \times \deg(v)\) matrix \(M(v)\) with complex entries \([M(v)]_{j_1,j_2}\), where \(j_1, j_2 \in \mathcal{S}(v)\) are edges incident with the vertex \(v\). The collection of such matrices for all \(v \in V\) will be denoted by \(\mathcal{M} = \{M(v)\}_{v \in V(G)}\).

Now to each non-trivial walk \(w = \{e', i_1, \ldots, i_N, e\}\) from \(e' \in E\) to \(e \in E\) on the graph \(G\) we associate a weight \(W(w)\) by

\[
W(w) = \prod_{k=1}^{\lfloor |w|_{comb} - 1 \rfloor} [M(v_k)]_{i_{k+1},i_k} \cdot [M(v_0)]_{i_1,e'},
\]

where \(v_0 = \partial(e')\), \(v_{|w|_{comb}} = \partial(e)\), \(v_k\) with \(k \in \{1, \ldots, |w|_{comb} - 1\}\) is the vertex incident with the internal line \(i_k\) as well as the internal line \(i_{k+1}\). To a trivial walk \(w = \{e', e\}\) we associate the weight

\[
W(w) = [M(\partial(e))]_{e,e'}.
\]

**Definition 2.3.** The generating function of walks from \(e' \in E\) to \(e \in E\) on the graph \(G\) associated with the collection \(\mathcal{M} = \{M(v)\}_{v \in V}\) is defined as

\[
T_{e,e'}(\beta) = \sum_{w \in \mathcal{W}_{e,e'}} W(w) e^{-\beta|w|} = \sum_{w \in \mathcal{W}_{e,e'}} W(w) e^{-\beta(\overline{w}(w), \overline{w})},
\]

where

\[
|w| = \langle \overline{w}(w), \overline{w} \rangle := \sum_{i \in \mathcal{I}} n_i(w) a_i.
\]

For given \(\mathcal{M}\) a walk \(w\) is called relevant if \(W(w) \neq 0\). The set of relevant walks from \(e'\) to \(e\) is denoted by \(\mathcal{W}_{e,e}(\mathcal{M})\).
Proposition 2.4. There is $\beta_0 > 0$ such that the series (2.3) converges for any $e, e' \in \mathcal{E}$ and all $\beta \in \mathbb{C}$ with $\Re \beta > \beta_0$. Moreover,

\[
\lim_{\Re \beta \to \infty} T_{e, e'}(\beta) = \begin{cases} 
[M(\partial(e))]_{e, e'} & \text{if } \partial(e) = \partial(e'), \\
0 & \text{otherwise}.
\end{cases}
\]

Definition 2.3 suggests that we write $W_{e, e'}$ as an infinite union of disjoint, non-empty sets by grouping together all walks $w$ with the same score $n(w)$,

\[
W_{e, e'}(n) = \{ w \in W_{e, e'} | n(w) = n \}
\]

such that

\[
W_{e, e'} = \bigcup_n W_{e, e'}(n).
\]

Note that these sets depend only on topology of the graph $G$ and are independent of its metric properties. Also if $w \in W_{e, e'}(n)$ then $w_{\text{rev}} \in W_{e', e}(n)$. $W_{e, e'}(\emptyset) = \emptyset$ if and only if $\partial(e) \neq \partial(e')$.

For the proof of Proposition 2.4 we need the following rather obvious fact:

Lemma 2.5. The sets $W_{e, e'}(n)$ are finite. Let

\[
|n| = \sum_{i \in I} n_i
\]

be the total number of internal lines traversed by any walk $w \in W_{e, e'}(n)$. The number of different walks in $W_{e, e'}(n)$ satisfies the bound

\[
|W_{e, e'}(n)| \leq \frac{|n|!}{\prod_{i \in I} n_i!}.
\]

Set

\[
T_{e, e'}(n) = \sum_{w \in W_{e, e'}(n)} W(w)
\]

if $W_{e, e'}(n)$ is nonempty and $T_{e, e'}(n) = 0$ whenever $W_{e, e'}(n) = \emptyset$. Observe that $T_{e, e'}(n)$ does not depend on the metric properties of the graph, i.e., is independent of the lengths of internal lines $\omega$.

For given $e, e' \in \mathcal{E}$ consider the set of scores of all walks from $e'$ to $e$,

\[
N_{e, e'} = \{ n | \text{there is a walk } w \in W_{e, e'}(n) \}.
\]

Since $n(w) = n(w_{\text{rev}})$, we have $N_{e, e'} = N_{e', e}$.

With this notation we have the following equivalent representation of (2.3):

\[
T_{e, e'}(\beta) = \sum_{n \in N_{e, e'}} T_{e, e'}(n) e^{-\beta(n, \omega)}.
\]

Obviously, the series in (2.3) converges absolutely if and only if the series in (2.3) does.

Proof of Proposition 2.4. Observe that

\[
\left| \sum_{w \in W_{e, e'}(n)} W(w) e^{-\beta(n, \omega)} \right| \leq \sum_{w \in W_{e, e'}(n)} \left( \max_{v \in V} \| M(v) \| \right)^{|n|+1} e^{-|n| \Re \beta a_{\min}},
\]
where
\begin{equation}
0 < a_{\min} := \min_{i \in I} a_i.
\end{equation}

From Lemma 2.5 and using the identity
\begin{equation}
\sum_{n \in \mathbb{N}^{|I|}} \prod_{i \in I} n_i! = |I|^N, \quad N \in \mathbb{N}
\end{equation}
we, therefore, obtain
\[
\left| \sum_{n \in \mathbb{N}^{|I|}} \left( \sum_{w \in \mathcal{W}_{e,e'}(n)} W(w) e^{-\beta \langle n, w \rangle} \right) \right| \leq \sum_{N=0}^{\infty} \left( \max_{v \in \mathcal{V}} \| M(v) \| \right)^{N+1} e^{-N \text{Re} \beta a_{\min} |I|^N}.
\]
This series converges for all $\beta \in \mathbb{C}$ with $\text{Re} \beta > \beta_0$, where
\begin{equation}
\beta_0 > \frac{1}{a_{\min}} \left( \max_{v \in \mathcal{V}} \log \| M(v) \| + \log |I| \right).
\end{equation}

We mention also the following simple result:

**Lemma 2.6 (Time Reversal Invariance).** If all matrices $M(v)$ are symmetric then so is the matrix $T(\beta)$ with matrix elements $T_{e,e'}(\beta)$ for all large $\text{Re} \beta > 0$. If all $M(v)$ are self-adjoint, then so is $T(\beta)$ for all large $\beta > 0$.

**Definition 2.7.** The family of matrices $\mathcal{M}$ is called combinatorial if every matrix entry of every matrix $M(v)$ equals either zero or one.

If $\mathcal{M}$ is combinatorial, the weight $W(w)$ of an arbitrary walk $w$ is either zero or one and we have the following simple result:

**Lemma 2.8.** If $\mathcal{M}$ is combinatorial and $\mathcal{W}_{e,e'}(\mathcal{M})$ finite then
\[
T_{e,e'}(0) = |\mathcal{W}_{e,e'}(\mathcal{M})|,
\]
i.e., the number of relevant walks from $e' \in \mathcal{E}$ to $e \in \mathcal{E}$.

We now provide some examples, which relate our formulation to well known combinatorial contexts. Viewing $\mathbb{Z}^2$ as a subset of $\mathbb{R}^2$, for an arbitrary $n \in \mathbb{N}$ consider the set
\[
\mathcal{V}_n = \left\{ (x_1, x_2) \in \mathbb{Z}^2 \mid 0 \leq x_2 \leq x_1 \leq n \right\}.
\]
We consider the non-compact graph $\mathcal{G}_n = (\mathcal{V}_n, \mathcal{I}, \mathcal{E}, \partial)$, where $\mathcal{E} = \{e, e'\}$, $\partial(e') = (0,0)$ and $\partial(e) = (n,n)$, and the vertices $v_1 \in \mathcal{V}_n$ and $v_2 \in \mathcal{V}_n$ are adjacent if and only if the Euclidean distance between these vertices is not larger than $\sqrt{2}$, $|v_1 - v_2| \leq \sqrt{2}$. Therefore, the set of internal lines $\mathcal{I}$ consists of all intervals joining the points of $\mathcal{V}_n$ and having Euclidean distance not greater than $\sqrt{2}$ (see Fig. 1). The metric distance between two adjacent vertices will be assumed to be equal 1, that is, $a_i = 1$ for all $i \in \mathcal{I}$.

**Example 2.9 (The Catalan numbers).** The number
\[
C_{n-1} = \frac{1}{n+1} \binom{2n}{n}, \quad n \in \mathbb{N}
\]
is called the \((n - 1)\)-th Catalan number (see, e.g., [29] and pp. 219 – 229 in [27]). Set \(K_{\text{Catalan}} = \{(1, 0), (0, 1)\}\). For an arbitrary vertex \(v \in V_n\) of the graph \(G_n\) and arbitrary \(j \in \mathcal{E} \cap \mathcal{I}\) adjacent to the vertex \(v\) we set

\[
\chi_v(j) = \begin{cases} 
  v' - v & \text{if } j \in \mathcal{I} \text{ and is adjacent to the vertex } v', \\
  (1, 0) \in \mathbb{Z}^2 & \text{if } j = e, \\
  (-1, 0) \in \mathbb{Z}^2 & \text{if } j = e'.
\end{cases}
\]

Let \(K_{\text{Schröder}} = \{(1, 0), (0, 1), (1, 1)\} \supset K_{\text{Catalan}}\) and

\[
[M^{\text{Schröder}}(v)]_{j_1, j_2} = \begin{cases} 
  1 & \text{if } \chi_v(j_1) \in K_{\text{Schröder}} \text{ and } -\chi_v(j_2) \in K_{\text{Schröder}}, \\
  0 & \text{otherwise}.
\end{cases}
\]

The set \(W_{e,e'}(M^{\text{Catalan}})\) is, obviously, finite. Therefore, the generating function \(T_{e,e'}(\beta)\) is entire. For given \(n \in \mathbb{N}\) the number \(T_{e,e'}(0)\) is the \((n - 1)\)-th Catalan number. The three other matrix elements \(T_{e,e}(\beta), T_{e',e}(\beta),\) and \(T_{e',e'}(\beta)\) vanish identically.

In the next example we continue with the same notation.

**Example 2.10 (The Schröder numbers).** The Schröder numbers (see, e.g., [29] and p. 178 in [27]) can be defined by the recurrence relation

\[
S_n = S_{n-1} + \sum_{k=0}^{n-1} S_k S_{n-k-1} \quad \text{with} \quad S_0 = 1.
\]

Let \(K_{\text{Schröder}} = \{(1, 0), (0, 1), (1, 1)\} \supset K_{\text{Catalan}}\) and

\[
[M^{\text{Schröder}}(v)]_{j_1, j_2} = \begin{cases} 
  1 & \text{if } \chi_v(j_1) \in K_{\text{Schröder}} \text{ and } -\chi_v(j_2) \in K_{\text{Schröder}}, \\
  0 & \text{otherwise}.
\end{cases}
\]

Obviously, \(W_{e,e'}(M^{\text{Catalan}}) \subseteq W_{e,e'}(M^{\text{Schröder}})\) is again a finite set. For given \(n \in \mathbb{N}\) the number \(T_{e,e'}(0)\) is now the \(n\)-th Schröder number. The three other matrix elements \(T_{e,e}(\beta), T_{e',e}(\beta),\) and \(T_{e',e'}(\beta)\) vanish identically.
For the next two examples consider the sets
\[ V_n^+ = \{(x_1, x_2) \in \mathbb{Z}^2 \mid 0 \leq x_1 \leq n, \ 0 \leq x_2 \leq n\}, \quad n \in \mathbb{N}. \]

Let \( G_n^+ = (V_n^+, \mathcal{I}, \mathcal{E}, \partial) \) be the non-compact graph with \( \mathcal{E} = \{e, e'\} \), \( \partial(e') = (n, 0) \), and the vertices \( v_1 \in V \) and \( v_2 \in V \) are adjacent if and only if the Euclidean distance between these vertices is not larger than \( \sqrt{2} \), \( |v_1 - v_2| \leq \sqrt{2} \). Therefore, the set of internal lines \( \mathcal{I} \) consists of all intervals joining the points of \( V_n^+ \) and having Euclidean length not greater than \( \sqrt{2} \) (see Fig. 2).

**Example 2.11 (Dyck paths).** Let \( K_{\text{Dyck}} = \{(1, 1), (1, -1)\} \) and
\[
[M_{\text{Dyck}}(v)]_{j_1, j_2} = \begin{cases} 
1 & \text{if } \chi_v(j_1) \in \mathcal{K}_{\text{Dyck}} \text{ and } -\chi_v(j_2) \in \mathcal{K}_{\text{Dyck}}, \\
0 & \text{otherwise}
\end{cases}
\]
if neither \( j_1 \) nor \( j_2 \) are external lines. We set
\[
[M_{\text{Dyck}}(v)]_{j_1, e'} = \begin{cases} 
1 & \text{if } \chi_v(j_1) \in \mathcal{K}_{\text{Dyck}}, \\
0 & \text{otherwise}
\end{cases}
\]
if \( j_2 = e' \) and
\[
[M_{\text{Dyck}}(v)]_{e, j_2} = \begin{cases} 
1 & \text{if } -\chi_v(j_2) \in \mathcal{K}_{\text{Dyck}}, \\
0 & \text{otherwise}
\end{cases}
\]
if \( j_1 = e. \)

Obviously, \( W_{e, e'}(M_{\text{Dyck}}) \) is a finite set. Therefore, \( T_{e, e'}(\beta) \) is entire, \( T_{e, e'}(0) \) is the number of Dyck paths on the graph \( G_n^+ \). A discussion of Dyck paths can be found in [11].

**Example 2.12 (Motzkin numbers).** The non-compact graphs \( G_n^+ \) are the same as for Dyck paths in Example 2.11. The Motzkin numbers (see, e.g., [11] and Problem 6.37 in [27]) can be defined by the recurrence relation
\[
M_n = M_{n-1} + \sum_{k=0}^{n-2} M_k M_{n-k-2} \quad \text{with} \quad M_0 = M_1 = 1.
\]
Set $K_{\text{Motzkin}} = \{(1, 1), (1, -1), (1, 0)\} \supset K_{\text{Dyck}}$ and

$$[M_{\text{Motzkin}}(v)]_{j_1, j_2} = \begin{cases} 1 & \text{if } \chi(v(j_1)) \in K_{\text{Motzkin}} \text{ and } \chi(v(j_2)) \in K_{\text{Motzkin}}, \\ 0 & \text{otherwise.} \end{cases}$$

Again the set $W_{e,e'}(M_{\text{Motzkin}})$ is finite and, therefore, $T_{e,e'}(\beta)$ is entire. For given $n \in \mathbb{N}$ the number $T_{e,e'}(0)$ is the $n$-th Motzkin number.

3. LAPLACE OPERATORS ON GRAPHS

In this section we will recall the theory of Laplace operators on a metric graph $G$ and the resulting scattering theory (see [12], [13], [14], [15], [16] for further details).

Given a finite non-compact graph $G = (V, I, E, \partial)$ with a metric structure $a = \{a_i\}_{i \in I}$ consider the Hilbert space

$$(3.1) \quad H \equiv H(I, E, a) = H_E \oplus H_I, \quad H_E = \bigoplus_{e \in E} \mathcal{H}_e, \quad H_I = \bigoplus_{i \in I} \mathcal{H}_i,$$

where $\mathcal{H}_e = L^2([0, \infty))$ for all $e \in E$ and $\mathcal{H}_i = L^2([0, a_i])$ for all $i \in I$. By $D_j$ with $j \in E \cup I$ denote the set of all $\psi_j \in H_j$ such that $\psi_j(x)$ and its derivative $\psi_j'(x)$ are absolutely continuous and $\psi_j''(x)$ is square integrable. Let $D^0_j$ denote the set of those elements $\psi_j$ of $D_j$ which satisfy

$$(3.2) \quad \psi_j(0) = \psi_j(a_j) = \psi_j'(0) = \psi_j'(a_j) = 0 \quad \text{for} \quad j \in I$$

and

$$\psi_j(0) = \psi_j'(0) = 0 \quad \text{for} \quad j \in E.$$

Let $\Delta^0$ be the differential operator

$$\Delta^0 \psi_j(x) = \frac{d^2}{dx^2} \psi_j(x), \quad j \in I \cup E$$

with $\psi = \{\psi_j\}_{j \in I \cup E}$ in the domain

$$\mathcal{D}^0 = \bigoplus_{j \in I \cup E} D^0_j \subset H.$$

It is straightforward to verify that $\Delta^0$ is a closed symmetric operator with deficiency indices equal to $|I| + 2|I|$.

We introduce an auxiliary finite-dimensional Hilbert space

$$(3.3) \quad K \equiv K(I, E) = K_E \oplus K_I^{-} \oplus K_I^+$$

with $K_E \cong \mathbb{C}^{|E|}$ and $K_I^{\pm} \cong \mathbb{C}^{|I|}$. The subspaces $K_I^{-}$ we associate with initial vertices of the internal lines $i \in I$, the subspaces $K_I^{+}$ with the terminal vertices. Let $dK$ denote the “double” of $K$, that is, $dK = K \oplus K$.

For any $\psi \in \mathcal{D} := \bigoplus_{j \in I \cup E} D_j$ we set

$$(3.4) \quad [\psi] := \underline{\psi} \oplus \underline{\psi}' \in dK,$$

with

$$\underline{\psi} = \left( \begin{array}{c} \{\psi_e(0)\}_{e \in E} \\
\{\psi_i(0)\}_{i \in I} \\
\{\psi_i(a_i)\}_{i \in I} \end{array} \right) \in K, \quad \underline{\psi}' = \left( \begin{array}{c} \{\psi_e'(0)\}_{e \in E} \\
\{\psi_i'(0)\}_{i \in I} \\
\{-\psi_i'(a_i)\}_{i \in I} \end{array} \right) \in K.$$

Here the vector notation is used with respect to the orthogonal decomposition $(3.3)$. 
To define the Laplace operator on the graph $G$ consider the family $\psi = \{\psi_j\}_{j \in \mathcal{E} \cup \mathcal{I}}$ of complex valued functions defined on $[0, \infty)$ if $j \in \mathcal{E}$ and on $[0, a_j]$ if $j \in \mathcal{I}$. Formally the (self-adjoint) Laplace operator is defined as

\begin{equation}
(\Delta(A, B, \alpha)\psi)_{j}(x) = \frac{d^2}{dx^2} \psi_j(x), \quad j \in \mathcal{I} \cup \mathcal{E}
\end{equation}

with the boundary conditions

\begin{equation}
A\psi + B\psi' = 0.
\end{equation}

By definition $A$ and $B$ are any complex $(|\mathcal{E}| + 2|\mathcal{I}|) \times (|\mathcal{E}| + 2|\mathcal{I}|)$ matrices such that

\begin{enumerate}[(i)]
  \item the matrix $(A, B)$ has maximal rank,
  \item the matrix $AB^\dagger$ is self-adjoint.
\end{enumerate}

Here and in what follows $(A, B)$ will denote the $(|\mathcal{E}| + 2|\mathcal{I}|) \times 2(|\mathcal{E}| + 2|\mathcal{I}|)$ matrix, where $A$ and $B$ are put next to each other.

The scattering matrix $S(k) = S(k; A, B, \alpha)$ associated to $\Delta(A, B, \alpha)$ has the following interpretation in terms of the solutions to the Schrödinger equation (see [12] and [13]). Consider the solutions $\psi^k_j(k \in \mathcal{E})$ of the stationary Schrödinger equation for $-\Delta(A, B, \alpha)$ at energy $k^2 > 0$,

\[-\Delta(A, B, \alpha)\psi^k_j(k) = k^2\psi^k_j(k)\]

of the form

\begin{equation}
\psi^k_j(x; k) = \begin{cases}
  S(k)_{jk} e^{ikx} & \text{for } j \in \mathcal{E}, j \neq k \\
  e^{-ikx} + S(k)_{kk} e^{ikx} & \text{for } j \in \mathcal{E}, j = k \\
  \alpha(k)_{jk} e^{ikx} + \beta(k)_{jk} e^{-ikx} & \text{for } j \in \mathcal{I}.
\end{cases}
\end{equation}

Thus, the number $S(k)_{jk}$ for $j \neq k$ is the transmission amplitude from channel $k \in \mathcal{E}$ to channel $j \in \mathcal{E}$ and $S(k)_{kk}$ is the reflection amplitude in channel $k \in \mathcal{E}$. Their absolute squares may be interpreted as transmission and reflection probabilities, respectively. The “interior” amplitudes

\[\alpha(k)_{jk} = \alpha(k; A, B, \alpha)_{jk}, \quad \beta(k)_{jk} = \beta(k; A, B, \alpha)_{jk}\]

are also of interest, since they describe how an incoming wave moves through a graph before it is scattered into an outgoing channel.

The condition for the $\psi^k_j(E)(k \in \mathcal{E})$ to satisfy the boundary conditions (5.7) immediately leads to the following solution for the scattering matrix $S(k): \mathcal{K}_\mathcal{E} \to \mathcal{K}_\mathcal{E}$ and the operators $\alpha(k)$ and $\beta(k)$ acting from $\mathcal{K}_\mathcal{E}$ to $\mathcal{K}_\mathcal{I}$. Indeed, by combining these operators into a map $\mathcal{K}_\mathcal{E}$ to $\mathcal{K} = \mathcal{K}_\mathcal{E} \oplus \mathcal{K}_\mathcal{I}^{(-)} \oplus \mathcal{K}_\mathcal{I}^{(+)}$ we obtain the linear equation

\begin{equation}
Z(k; A, B, \alpha) \begin{pmatrix} S(k) \\ \alpha(k) \\ \beta(k) \end{pmatrix} = -(A - ikB) \begin{pmatrix} \mathbb{1} \\ 0 \\ 0 \end{pmatrix}
\end{equation}

with

\begin{equation}
Z(k; A, B, \alpha) = AX(k; \alpha) + ikBY(k; \alpha),
\end{equation}

where

\begin{equation}
X(k; \alpha) = \begin{pmatrix} \mathbb{1} & 0 & 0 \\ 0 & \mathbb{1} & \mathbb{1} \\ 0 & e^{ik\alpha} & e^{-ik\alpha} \end{pmatrix}, \quad Y(k; \alpha) = \begin{pmatrix} \mathbb{1} & 0 & 0 \\ 0 & \mathbb{1} & -\mathbb{1} \\ 0 & -e^{ik\alpha} & e^{-ik\alpha} \end{pmatrix}.
\end{equation}
The diagonal $|I| \times |I|$ matrices $e^{\pm ik}a$ are given by
\begin{equation}
(3.13) \quad e^{\pm ik}a_{jk} = \delta_{jk} e^{\pm ikj} \quad \text{for} \ j, k \in I.
\end{equation}

**Theorem 3.1** (= Theorem 3.2 in \cite{12}). For any $k \in \mathbb{R}$
\begin{equation}
(3.10) \quad \text{Ran} \left( A - ikB \right) \begin{pmatrix} \mathbb{I} \\ 0 \\ 0 \end{pmatrix} \subset \text{Ran} \left( Z(k; A, B, a) \right).
\end{equation}
Thus, equation (3.10) has a solution even if $\det Z(k; A, B, a) = 0$ for some $k \in \mathbb{R}$. This solution defines the scattering matrix uniquely. Moreover,
\begin{equation}
(3.14) \quad S(k) = - \begin{pmatrix} 0 \end{pmatrix} Z(k; A, B, a)^{-1} P_{\ker Z(k; A, B, a)} (A - ikB) \begin{pmatrix} \mathbb{I} \\ 0 \\ 0 \end{pmatrix}
\end{equation}
is unitary for all $k \in \mathbb{R} \setminus \{0\}$.

In the case with no internal lines ($I = \emptyset$) the relation (3.14) for the scattering matrix simplifies to
\begin{equation}
(3.15) \quad S(k; A, B) = - (A + ikB)^{-1} (A - ikB).
\end{equation}

**Proposition 3.2.** If $\det (A + ikB) = 0$ for some $k \in \mathbb{C}$, then $k = \kappa$ with $\kappa \in \mathbb{R}$. For any sufficiently large $\rho > 0$ there is a constant $C_\rho > 0$ such that
\begin{equation}
(3.16) \quad \|(A + ikB)^{-1}\| \leq C_\rho (1 + |k|)^{-1}
\end{equation}
for all $k \in \mathbb{C}$ with $|k| > \rho$.

**Proof.** Assume that $\det (A + ikB) = 0$ for some $k \in \mathbb{C}$ with $\text{Re} \, k \neq 0$. Then also
\begin{equation}
\det (A^\dagger - i\kappa B^\dagger) = \overline{\det (A + ikB)} = 0.
\end{equation}
Therefore, there is a $\chi \neq 0$ such that
\begin{equation}
(3.17) \quad (A^\dagger - i\kappa B^\dagger) \chi = 0.
\end{equation}
In particular, we have $(BA^\dagger - i\kappa BB^\dagger) \chi = 0$. Therefore, since $BA^\dagger$ is self-adjoint, we get
\begin{align*}
\langle \chi, BA^\dagger \chi \rangle &= \langle \chi, BB^\dagger \chi \rangle \text{Im} \, k, \\
\langle \chi, BB^\dagger \chi \rangle \text{Re} \, k &= 0.
\end{align*}
The second equality implies that $\chi \in \ker B^\dagger$. Then, by (3.17), $\chi \in \ker A^\dagger$. Since the matrix $(A, B)$ is of maximal rank, we have $\ker A^\dagger \cap \ker B^\dagger = \{0\}$. Thus, $\chi = 0$ which contradicts the assumption and, hence, $\text{Re} \, k = 0$.

Since $\det (A + ikB)$ is a polynomial in $k$, it has a finite number of zeroes. Take an arbitrary $\rho > 0$ such that all its zeroes lie in the disk $|k| < \rho$. Using the matrix inverse formula we represent any element of $(A + ikB)^{-1}$ as a quotient of two polynomials of degrees $|E| + 2|I| - 1$ and $|E| + 2|I|$, respectively. In turn, this implies the estimate (3.16). \hfill \Box

**Theorem 3.3.** The scattering matrix $S(k) = S(k; A, B, a)$ is a meromorphic function in the complex $k$-plane. In upper half-plane $\text{Im} \, k > 0$ it has at most a finite number of poles.
which are located on the imaginary semiaxis \( \text{Re} \, k = 0 \). Outside these poles the scattering matrix is holomorphic for all \( \text{Im} \, k > 0 \) and determined by the relation
\[
(3.18) \quad S(k) = -\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} Z(k; A, B, \underline{a})^{-1} (A - ikB) \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}.
\]

**Proof.** Assume that \( \det Z(k; A, B, \underline{a}) = 0 \) for some \( k \in \mathbb{C} \) with \( \text{Im} \, k > 0 \) and \( \text{Re} \, k \neq 0 \). This implies that the homogeneous equation
\[
Z(k; A, B, \underline{a}) \begin{pmatrix} s \\ \alpha \\ \beta \end{pmatrix} = 0
\]
has a nontrivial solution with \( s \in \mathcal{K}_E \) and \( \alpha, \beta \in \mathbb{C}[\mathbb{I}] \). Consider the function \( \psi(x) = \{\psi_j(x)\}_{j \in \mathbb{I} \cup \mathcal{E}} \) defined by
\[
\psi_j(x) = \begin{cases} s_j e^{ikx} & \text{for } j \in \mathcal{E}, \\ \alpha_j e^{ikx} + \beta_j e^{-ikx} & \text{for } j \in \mathbb{I}. \end{cases}
\]
Obviously, \( \psi(x) \) satisfies the boundary conditions (3.7). Moreover, \( \psi \in L^2(\mathcal{G}) \) since \( \text{Im} \, k > 0 \). Hence, \( k^2 \in \mathbb{C} \) with \( \text{Im} \, k^2 \neq 0 \) is an eigenvalue of the operator \( \Delta(A, B, \underline{a}) \) which contradicts the self-adjointness of \( \Delta(A, B, \underline{a}) \).

Since \( \det Z(k; A, B, \underline{a}) \) is an entire function in \( k \) which does not vanish identically, from (3.10) it follows that the scattering matrix \( S(k) \) is a meromorphic function in the complex \( k \)-plane. To prove that the scattering matrix \( S(k) \) has at most a finite number of poles on the imaginary semiaxis \( \{k \in \mathbb{C} \mid \text{Re} \, k = 0, \, \text{Im} \, k > 0\} \) it suffices to show that the determinant \( \det Z(k; A, B, \underline{a}) \) does not vanish for all sufficiently large \( \text{Im} \, k > 0 \). To see this we set \( k = i\kappa \) with \( \kappa > 0 \) and assume there is an unbounded non-decreasing sequence \( \{\kappa_k\}_{k \in \mathbb{N}} \) such that
\[
\det Z(i\kappa_k; A, B, \underline{a}) = 0 \quad \text{for all} \quad k \in \mathbb{N}.
\]
Therefore, there is a sequence \( \{\chi_k\}_{k \in \mathbb{N}} \) of normalized elements \( \chi_k \in \mathcal{K} \) such that
\[
X(i\kappa_k; \underline{a})^\dagger A^\dagger \chi_k = \kappa_k Y(i\kappa_k; \underline{a})^\dagger B^\dagger \chi_k.
\]
It is straightforward to verify that \( X(i\kappa_k; \underline{a}) \) is invertible and
\[
R_k := \left( X(i\kappa_k; \underline{a})^\dagger \right)^{-1} Y(i\kappa_k; \underline{a})^\dagger = \begin{pmatrix} 0 & 0 \\ 0 & \coth(\kappa \underline{a}) \end{pmatrix}^{-1} - \begin{pmatrix} 0 \\ -[\sinh(\kappa \underline{a})]^{-1} \coth(\kappa \underline{a}) \end{pmatrix}
\]
with a notation analogous to (3.13). Thus,
\[
(3.19) \quad (A^\dagger - \kappa_k B^\dagger) \chi_k = \kappa_k (R_k - I) B^\dagger \chi_k
\]
for all \( k \in \mathbb{N} \). Observe that \( \|R_k - I\| = O(e^{-c_\kappa k}) \) for some \( c > 0 \) as \( k \to \infty \). By Proposition 5.2 the operator \( A^\dagger - \kappa B^\dagger \) is invertible for all sufficiently large \( \kappa \). Moreover, \( \|(A^\dagger - \kappa B^\dagger)^{-1}\| \leq C \) with \( C > 0 \) for all sufficiently large \( \kappa \). Thus, equation (3.19) implies that \( \chi_k \to 0 \) which contradicts the assumption \( \|\chi_k\| = 1 \). \( \square \)

In the lower half-plane \( \text{Im} \, k < 0 \) the scattering matrix may have poles with \( \text{Re} \, k \neq 0 \) (see, e.g., Example 3.2 in [12]). These poles correspond to resonances.

The notion of local boundary conditions has been introduced in our article [12] and is discussed in more details in [15] and [16]. Local boundary conditions couple only
those boundary values of $\psi$ and of its derivative $\psi'$ which belong to the same vertex. The precise definition is as follows.

With respect to the orthogonal decomposition $\mathcal{K} = \mathcal{K}_E \oplus \mathcal{K}_I^(-) \oplus \mathcal{K}_I^+$ any element $z$ of $\mathcal{K}$ can be represented as a vector

\[
 z = \begin{pmatrix}
 \{ z_e \}_{e \in \mathcal{E}} \\
 \{ z_i^(-) \}_{i \in \mathcal{I}} \\
 \{ z_i^+ \}_{i \in \mathcal{I}}
\end{pmatrix}.
\]

Consider the orthogonal decomposition

\[
 \mathcal{K} = \bigoplus_{v \in V} \mathcal{L}_v
\]

with $\mathcal{L}_v$ being the linear subspace of dimension $\deg(v)$ spanned by those elements (3.20) of $\mathcal{K}$ which satisfy

\[
 z_e = 0 \quad \text{if} \quad e \in \mathcal{E} \quad \text{is not incident with the vertex} \quad v,
\]

\[
 z_i^(-) = 0 \quad \text{if} \quad v \quad \text{is not an initial vertex of} \quad i \in \mathcal{I},
\]

\[
 z_i^+ = 0 \quad \text{if} \quad v \quad \text{is not a terminal vertex of} \quad i \in \mathcal{I}.
\]

Set $d\mathcal{L}_v := \mathcal{L}_v \oplus \mathcal{L}_v \cong \mathbb{C}^{2\deg(v)}$. By the First Theorem of Graph Theory we have

\[
 \sum_{v \in V(\mathcal{G})} \deg(v) = |\mathcal{E}| + 2|\mathcal{I}|
\]

such that

\[
 \bigoplus_{v \in V(\mathcal{G})} d\mathcal{L}_v = \mathcal{K}.
\]

**Definition 3.4.** Given the graph $\mathcal{G} = \mathcal{G}(V, \mathcal{I}, \mathcal{E}, \partial)$, the boundary conditions $(A, B)$ satisfying (3.8) are called local on $\mathcal{G}$ if and only if there is an invertible map $C : \mathcal{K} \to \mathcal{K}$ and linear transformations $A(v)$ and $B(v)$ in $\mathcal{L}_v$ such that the direct sum decompositions

\[
 CA = \bigoplus_{v \in V} A(v) \quad \text{and} \quad CB = \bigoplus_{v \in V} B(v)
\]

hold simultaneously. Otherwise the boundary conditions are called non-local.

For instance, for a single-vertex graph any boundary conditions are local. The boundary conditions considered in Example 3.4 of [16] are non-local.

**4. Combinatorial Fourier Expansion of the Scattering Matrix**

In this section we will perform a harmonic analysis of the scattering matrix with respect to the lengths $\mathbf{a} = \{ a_i \}_{i \in \mathcal{I}} \in (\mathbb{R}^+)^{\mathcal{I}}$ of the internal lines of the graph $\mathcal{G}$. The main results of this section are presented in Theorems 4.2 and 4.10. In Theorem 4.2 the absolute convergence of the Fourier series for the scattering matrix is proved. Theorem 4.10 expresses its Fourier coefficients as sums over the walks on the graph. Combining these two results proves the combinatorial Fourier expansion formula (4.18).

Throughout the whole section we will assume that the (topological) graph $\mathcal{G}$ as well as the boundary conditions $(A, B)$ are fixed. To carry out the analysis we will now treat $\mathbf{a}$ as a parameter which may belong to $\mathbb{R}^{\mathcal{I}}$ or even $\mathbb{C}^{\mathcal{I}}$.

We start with the following simple but important observation.
Lemma 4.1. For arbitrary \( k > 0 \) the scattering matrix \( S(k; A, B, \underline{a}) \) is uniquely defined as a solution of (3.10) for all \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \). Moreover, the scattering matrix is periodic with respect to \( \underline{a} \),

\[
S \left( k; A, B, \underline{a} + \frac{2\pi k}{k} \ell \right) = S(k; A, B, \underline{a})
\]

for arbitrary \( \ell \in \mathbb{Z}^{|\mathcal{I}|} \).

Proof. It suffices to consider those \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \) for which \( \det Z(k; A, B, \underline{a}) = 0 \), since the claim is obvious when the determinant is non-vanishing. For \( \underline{a} \in (\mathbb{R}_+)^{|\mathcal{I}|} \) the fact that \( S(k; A, B, \underline{a}) \) is uniquely defined as a solution of (3.10) is guaranteed by Theorem 3.1.

The case of arbitrary \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \) can be treated exactly in the same way (see the proof of Theorem 3.2 in [12]).

The periodicity follows immediately from (3.10) and the fact that the matrices \( X(k; \underline{a}) \) and \( Y(k; \underline{a}) \) in (3.12) are \( \frac{2\pi k}{k} Z^{|\mathcal{I}|} \)-periodic.

Lemma 4.1 suggests to consider a Fourier expansion of the scattering matrix. The following theorem ensures the absolute convergence of the corresponding Fourier series.

Theorem 4.2. Let \( k > 0 \) be arbitrary. For all \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \) the Fourier expansion of the scattering matrix

\[
S(k; A, B, \underline{a}) = \sum_{\underline{n} \in \mathbb{Z}^{|\mathcal{I}|}} \hat{S}_n(k; A, B) e^{i k \langle \underline{n}, \underline{a} \rangle}
\]

with

\[
\hat{S}_n(k; A, B) = \left( \frac{k}{2\pi} \right)^{|\mathcal{I}|} \int_{[0, 2\pi/k]^{|\mathcal{I}|}} d\underline{a} S(k; A, B, \underline{a}) e^{-i k \langle \underline{n}, \underline{a} \rangle}
\]

converges absolutely and uniformly on compact subsets of \( \mathbb{R}^{|\mathcal{I}|} \). The Fourier coefficients \( \hat{S}_n \) vanish for all \( \underline{n} = \{n_i\}_{i \in \mathcal{I}} \in \mathbb{Z}^{|\mathcal{I}|} \) with \( n_i < 0 \) for at least one \( i \in \mathcal{I} \).

For the proof we need a couple of auxiliary results. Set

\[
\mathcal{A} = \{ \underline{a} = \{a_i\}_{i \in \mathcal{I}} \mid \Re a_i \in \mathbb{R}, \Im a_i > 0 \} \subset \mathbb{C}^{|\mathcal{I}|}.
\]

Lemma 4.3. For any \( k > 0 \) the determinant \( \det Z(k; A, B, \underline{a}) \) has no zeroes for all \( \underline{a} \in \mathcal{A} \).

Proof. Assume there is \( \underline{a} \in \mathcal{A} \) such that \( \det Z(k; A, B, \underline{a}) = 0 \). Then there are \( s \in \mathbb{C}^{|\mathcal{E}|} \) and \( \alpha, \beta \in \mathbb{C}^{|\mathcal{I}|} \) such that

\[
Z(k; A, B, \underline{a}) \begin{pmatrix} s \\ \alpha \\ \beta \end{pmatrix} = 0.
\]

Equivalently this gives

\[
(A + i k B) \begin{pmatrix} s \\ \alpha \\ e^{-i k a_\beta} \beta \end{pmatrix} + (A - i k B) \begin{pmatrix} 0 \\ \beta \\ e^{i k a_\alpha} \alpha \end{pmatrix} = 0.
\]

The operator \( (A + i k B)^{-1} (A - i k B) \) is unitary for all \( k > 0 \) (see the proof of Theorem 2.1 in [12]). Since unitary transformations preserve the canonical Hilbert norm on \( \mathbb{C}^{|\mathcal{E}|} + 2 |\mathcal{I}| \),
we have
\[ \|s\|^2 + \sum_{i \in \mathcal{I}} |\alpha_i|^2 (1 - e^{-2k\text{Im} \alpha_i}) + \sum_{i \in \mathcal{I}} |\beta_i|^2 (e^{2k\text{Im} \alpha_i} - 1) = 0, \]
which implies \( s = 0 \) and \( \alpha = \beta = 0 \). \( \square \)

**Proposition 4.4.** Let \( k > 0 \) be arbitrary. For all \( \underline{a} \in \text{clos}(\mathcal{A}) := \{ \underline{a} = \{a_i\}_{i \in \mathcal{I}} | \text{Re } a_i \in \mathbb{R}, \text{Im } a_i \geq 0 \} \) the scattering matrix \( S(k; A, B, \underline{a}) \) is uniquely defined as a solution of (3.10) and satisfies the bound
\[ (4.3) \quad \|S(k; A, B, \underline{a})\| \leq 1. \]
Moreover, it is a rational function of \( \ell = \{t_i\}_{i \in \mathcal{I}} \) with \( t_i := e^{ik\alpha_i} \), i.e. a quotient of \( \mathcal{B}(\mathcal{K}_\mathcal{I}) \)-valued polynomials in the variables \( t_i \). Thus, for all \( \underline{a} \in \text{clos}(\mathcal{A}) \) the scattering matrix is \( \frac{2\pi}{k} \mathbb{Z}^{|\mathcal{I}|} \)-periodic,
\[ S \left( k; A, B, \underline{a} + \frac{2\pi}{k} \ell \right) = S(k, A, B, \underline{a}), \quad \ell \in \mathbb{Z}^{|\mathcal{I}|}. \]

**Proof.** By Lemma 4.3 equation (3.10) has a unique solution for all \( \underline{a} \in \mathcal{A} \). Equations (3.11) and (3.12) imply that \( Z(k; A, B, \underline{a}) \) is a polynomial function of the components of \( \ell \). Obviously, \( Z(k; A, B, \underline{a})^{-1} \) is also a rational function of \( \ell \). Thus, by (3.10) the scattering matrix \( S(k; A, B, \underline{a}) \) is a rational function of \( \ell \). Thus, it is \( \frac{2\pi}{k} \mathbb{Z}^{|\mathcal{I}|} \)-periodic.

Using (3.10) it is easy to check that this solution satisfies the relation
\[ \begin{pmatrix} S(k; A, B, \underline{a}) \\ e^{-ik\ell} \beta(k; A, B, \underline{a}) \end{pmatrix} = -(A + ikB)^{-1}(A - ikB) \begin{pmatrix} \mathbb{I} \\ \beta(k; A, B, \underline{a}) \end{pmatrix}. \]
Since \( (A + ikB)^{-1}(A - ikB) \) is unitary we obtain
\[ S(k; A, B, \underline{a})^\dagger S(k; A, B, \underline{a}) + \alpha(k; A, B, \underline{a})^\dagger (\mathbb{I} - e^{-2k\text{Im} \underline{a}}) \alpha(k; A, B, \underline{a}) \]
\[ + \beta(k; A, B, \underline{a})^\dagger (e^{2k\text{Im} \underline{a}} - \mathbb{I}) \beta(k; A, B, \underline{a}) = \mathbb{I}, \]
where \( \text{Im } \underline{a} = \{ \text{Im } a_i \}_{i \in \mathcal{I}} \). From (4.4) it follows immediately that
\[ 0 \leq S(k; A, B, \underline{a})^\dagger S(k; A, B, \underline{a}) \leq \mathbb{I} \]
in the operator sense. This proves the bound (4.3) for all \( \underline{a} \in \mathcal{A} \). Recalling Lemma 4.1 completes the proof. \( \square \)

A priori it is not clear whether the boundary values of the scattering matrix \( S(k; A, B, \underline{a}) \) with \( \underline{a} \in \mathcal{A} \) coincide with those given by equation (3.14) for all \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \). The following lemma shows the “non-tangential continuity” of the scattering matrix \( S(k; A, B, \underline{a}) \) with respect to \( \underline{a} \in \text{clos}(\mathcal{A}) \).

**Lemma 4.5.** Let \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \) and \( k > 0 \) be arbitrary. For any sequence \( \{a_j\}_{j \in \mathbb{N}}, \underline{a}_j \in \mathcal{A} \) converging to \( \underline{a} \in \mathbb{R}^{|\mathcal{I}|} \) the relation
\[ (4.5) \quad \lim_{j \to \infty} S(k; A, B, \underline{a}_j) = S(k; A, B, \underline{a}) \]
holds.

For the proof we need the following elementary result.
Lemma 4.6. Let \( T_n \) be a sequence of invertible operators on the finite-dimensional Hilbert space \( \mathcal{H} \) converging to the operator \( T \). Then
\[
\lim_{n \to \infty} P_{\text{Ker} T} T_n^{-1} P_{\text{Ker} T}^\dagger = P_{\text{Ker} T} T^{-1} P_{\text{Ker} T}^\dagger,
\]
where \( P_{\mathcal{K}}^\perp \) denotes the orthogonal projection onto orthogonal complement in \( \mathcal{H} \) of the subspace \( \mathcal{K} \subseteq \mathcal{H} \).

Proof. Consider the operators \( T_n \) and \( T \) as maps from \( (\text{Ker} T)^\perp \) to \( (\text{Ker} T)^\perp \). Since these maps are invertible, the claim follows from the obvious relation
\[
T_n^{-1} = T^{-1} \left[ I + T^{-1}(T_n - T) \right]^{-1}.
\]

Proof of Lemma 4.6. Introduce the shorthand notation
\[
Z(\bar{a}) \equiv Z(k; A, B, \bar{a}) \quad \text{and} \quad S(\bar{a}) \equiv S(k; A, B, \bar{a}).
\]
From Theorem 3.1 and Lemma 4.3 it follows that
\[
S(\bar{a}) = - \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a})^{-1} P_{\text{Ker} Z(\bar{a})}^\dagger (A - ikB) \left( \begin{array}{c} 1 \\ 0 \end{array} \right)
\]
and
\[
S(\bar{a}_j) = - \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a}_j)^{-1} (A - ikB) \left( \begin{array}{c} 1 \\ 0 \end{array} \right)
\]
\[
= - \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a}_j)^{-1} P_{\text{Ker} Z(\bar{a}_j)}^\dagger (A - ikB) \left( \begin{array}{c} 1 \\ 0 \end{array} \right).
\]
Thus, to prove the claim it suffices to show that
\[
\lim_{j \to \infty} \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a}_j)^{-1} P_{\text{Ker} Z(\bar{a}_j)}^\dagger = \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a})^{-1} P_{\text{Ker} Z(\bar{a})}^\dagger.
\]

From Theorem 3.1 in [12] it follows that all elements \( z \) of \( Z(\bar{a}) \) satisfy \( P_{\mathcal{K}} z = 0 \), where \( P_{\mathcal{K}} \) is the orthogonal projection in \( \mathcal{K} \) onto \( \mathcal{K} \). Thus,
\[
\left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a}_j)^{-1} P_{\text{Ker} Z(\bar{a}_j)}^\dagger = \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) P_{\text{Ker} Z(\bar{a}_j)} Z(\bar{a}_j)^{-1} P_{\text{Ker} Z(\bar{a}_j)}^\dagger
\]
for any \( j \in \mathbb{N} \) and
\[
\left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) Z(\bar{a})^{-1} P_{\text{Ker} Z(\bar{a})}^\dagger = \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) P_{\text{Ker} Z(\bar{a})} Z(\bar{a})^{-1} P_{\text{Ker} Z(\bar{a})}^\dagger.
\]
By Lemma 4.6 we have
\[
\lim_{j \to \infty} P_{\text{Ker} Z(\bar{a}_j)} Z(\bar{a}_j)^{-1} P_{\text{Ker} Z(\bar{a}_j)}^\dagger = P_{\text{Ker} Z(\bar{a})} Z(\bar{a})^{-1} P_{\text{Ker} Z(\bar{a})}^\dagger.
\]
Combining this with (4.7) and (4.8) we obtain (4.6).

For fixed \( k > 0 \), consider
\[
F(t) := S(k; A, B, \bar{a}) \quad \text{with} \quad t = e^{ik}.
\]
Recall that by (3.10) – (3.12) the scattering matrix \( S(k; A, B, \bar{a}) \) depends on \( \bar{a} \) only through \( t = e^{ik} \). The map \( \bar{a} \mapsto e^{ik} \) maps the set
\[
\left\{ \bar{a} \in \mathbb{C}^{|\mathcal{I}|} | \bar{a} = \{a_i\}_{i \in \mathcal{I}} \right\}
\]
with \( 0 < \text{Re} a_i \leq 2\pi/k \) and \( \text{Im} a_i > 0 \) for all \( i \in \mathcal{I} \).
proof of Theorem 4.2. Since by Proposition 4.4, \( F(\xi) \) is a rational \( B(\mathcal{K}_E) \)-valued function, it can be analytically continued as a meromorphic function on all of \( \xi \in \mathbb{C}^{\mathbb{Z}} \). Moreover, it is holomorphic in the polydisc \( \mathbb{D}^{\mid \mathbb{Z}} \setminus \{0\} \). By (4.3) we have \( \|F(\xi)\| \leq 1 \) for all \( \xi \in \mathbb{D}^{\mid \mathbb{Z}} \setminus \{0\} \). Therefore, the Laurent expansion of \( F \) contains no terms with negative powers. Thus, \( F \) is holomorphic in \( \mathbb{D}^{\mid \mathbb{Z}} \).

The bound (4.3) also implies that

\[
\sup_{r \in [0,1)} \int_{\mathbb{T}^{\mid \mathbb{Z}}} \|F(r \xi)\|^p d\mu(\xi) \leq \mu(\mathbb{T}^{\mid \mathbb{Z}})
\]

for any \( p \in (0, \infty) \), where \( \mu \) stands for the Haar measure on the torus \( \mathbb{T}^{\mid \mathbb{Z}} \) and

\[
\sup_{r \in [0,1)} \sup_{\xi \in \mathbb{T}^{\mid \mathbb{Z}}} \|F(r \xi)\| \leq 1.
\]

For every \( \xi \in \mathbb{T}^{\mid \mathbb{Z}} \) the operator \( F(\xi) \) is unitary, which means that \( F \) is an inner function.

**Proof of Theorem 4.2.** Since by Proposition 4.4, \( F(\xi) \) is a rational \( B(\mathcal{K}_E) \)-valued function, it can be analytically continued as a meromorphic function on all of \( \xi \in \mathbb{C}^{\mathbb{Z}} \). Moreover, it is holomorphic in the polydisc \( \mathbb{D}^{\mid \mathbb{Z}} \setminus \{0\} \). By (4.3) we have \( \|F(\xi)\| \leq 1 \) for all \( \xi \in \mathbb{D}^{\mid \mathbb{Z}} \setminus \{0\} \). Therefore, the Laurent expansion of \( F \) contains no terms with negative powers. Thus, \( F \) is holomorphic in \( \mathbb{D}^{\mid \mathbb{Z}} \).

The bound (4.3) also implies that

\[
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\]

for any \( p \in (0, \infty) \), where \( \mu \) stands for the Haar measure on the torus \( \mathbb{T}^{\mid \mathbb{Z}} \) and

\[
\sup_{r \in [0,1)} \sup_{\xi \in \mathbb{T}^{\mid \mathbb{Z}}} \|F(r \xi)\| \leq 1.
\]

For every \( \xi \in \mathbb{T}^{\mid \mathbb{Z}} \) the operator \( F(\xi) \) is unitary, which means that \( F \) is an inner function. \( \square \)

**Definition 4.4.** Let \( G = (V, I, E, \partial) \) be a graph of arbitrary boundary conditions. We associate the single-vertex graph \( G_v = (\{v\}, I_v, E_v, \partial_v) \) with the following properties

(i) \( I_v = \emptyset \),
(ii) \( \partial_v(e) = v \) for all \( e \in E_v \),
(iii) \( |E_v| = \deg_G(v) \), the degree of the vertex \( v \) in the graph \( G \),
(iv) there is an injective map \( \Psi_v : E_v \to \mathcal{E} \cup \mathcal{I} \) such that \( v \in \partial \circ \Psi_v(e) \) for all \( e \in E_v \).

Since the boundary conditions are assumed to be local (see Definition 3.13), we can consider the Laplace operator \( \Delta(A_v, B_v) \) on \( L^2(G_v) \) associated with the boundary conditions \((A_v, B_v)\) induced by \((A, B)\), see (3.22). By (3.15) the scattering matrix for \( \Delta(A_v, B_v) \) is given by

\[
S_v(k) = -(A_v + ikB_v)^{-1}(A_v - ikB_v).
\]
Now to each walk \( w = \{e', i_1, \ldots, i_N, e\} \) from \( e' \in \mathcal{E} \) to \( e \in \mathcal{E} \) on the graph \( \mathcal{G} \) similar to \([2.1]\) we associate a weight \( W(w; k) \) by
\[
W(w; k) = e^{ik\langle w; \varrho \rangle} \tilde{W}(w; k)
\]
with
\[
\tilde{W}(w; k) = \prod_{k=0}^{w_{\text{comb}}} S_{\nu_k}(k) e_k^{(+)} e_k^{(-)}.
\]
Here \( e_k^{(\pm)} \in \mathcal{E}_{\nu_k} \) are defined as
\[
e_k^{(-)} = \begin{cases} 
\Psi_{\nu_k}^{-1}(i_k), & \text{if } 1 \leq k \leq |w|_{\text{comb}}, \\
\Psi_{\nu_k}^{-1}(e), & \text{if } k = 0,
\end{cases}
\]
and
\[
e_k^{(+)} = \begin{cases} 
\Psi_{\nu_k}^{-1}(i_{k+1}), & \text{if } 0 \leq k \leq |w|_{\text{comb}} - 1, \\
\Psi_{\nu_k}^{-1}(e'), & \text{if } k = |w|_{\text{comb}} + 1,
\end{cases}
\]
where the map \( \Psi_v \) is defined by Definition \([4.9]\). Note that \( \tilde{W}(w; k) \) is independent of the metric properties of the graph. Obviously, for a trivial walk \( w = \{e', e\} \) we have \( \tilde{W}(w; k) = S_{\nu}(k) \Psi_v^{-1}(e), \Psi_v^{-1}(e') \), where \( v = \partial(e) = \partial(e') \).

**Theorem 4.10.** The matrix elements of the \( \mathbf{m} \)-th Fourier coefficients \([4.2]\) are given by the sum over the walks with score \( \mathbf{n} \)
\[
[S_{\mathbf{m}}(k; A, B)]_{e, e'} = \sum_{w \in \mathcal{W}_{e, e'} (\mathbf{n})} \tilde{W}(w; k)
\]
if \( \mathcal{W}_{e, e'} (\mathbf{n}) \) is nonempty and \( [S_{\mathbf{m}}(k; A, B)]_{e, e'} = 0 \) whenever \( \mathcal{W}_{e, e'} (\mathbf{n}) = \emptyset \).

**Proof.** Obviously, it suffices to show that the \( \mathbf{m} \)-th coefficient of the multi-dimensional Taylor expansion of the scattering matrix \( S(k; A, B, \varrho) \) with respect to \( \{t_i\} \in \mathbb{D}^{[\mathcal{I}]} \) with \( t_i := e^{ik_{a_i}} \) coincides with the r.h.s. of \([4.12]\). Recall that by Theorem \([3.1]\) and Lemma \([4.3]\) for all \( \varrho \in A \) the scattering matrix is given by
\[
S(k; A, B, \varrho) = - (I \ 0 \ 0) (AX(k; \varrho) + 1kBY(k; \varrho))^{-1} (A - 1kB) \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix},
\]
where \( X(k; \varrho) \) and \( Y(k; \varrho) \) were defined in \([3.12]\). Obviously,
\[
AX(k; \varrho) + 1kBY(k; \varrho) = (A + 1kB)U(k; \varrho) + (A - 1kB)R(k; \varrho)
\]
\[
= (A + 1kB)[I + (A + 1kB)^{-1}(A - 1kB)R(k; \varrho)U(k; \varrho)^{-1}]U(k; \varrho),
\]
where
\[
U(k; \varrho) := \begin{pmatrix} I & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & e^{-ik_{\varrho}} \end{pmatrix}
\]
and
\[
R(k; \varrho) := X(k; \varrho) - U(k; \varrho) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & I \\ 0 & e^{ik_{\varrho}} & 0 \end{pmatrix}
\]
with respect to the orthogonal decomposition (3.3). Equation (4.14) implies that

\[(AX(k; \underline{a}) + i k BY(k; \underline{a}))^{-1} = U(k; \underline{a})^{-1} \sum_{n=0}^{\infty} [- (A + i k B)^{-1} (A - i k B)GH(k; \underline{a})]^n (A + i k B)^{-1}\]

with

(4.16) \[G = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & I \\ 0 & I & 0 \end{pmatrix} \quad \text{and} \quad H(k; \underline{a}) = \begin{pmatrix} I & 0 & 0 \\ 0 & e^{ik\underline{a}} & 0 \\ 0 & 0 & e^{i k \underline{a}} \end{pmatrix}\]

such that \( R(k; \underline{a})U(k; \underline{a})^{-1} = GH(k; \underline{a}) \). Combining this representation with (4.13) we obtain

\[S(k; A, B, \underline{a}) = \sum_{n=0}^{\infty} (I \ 0 \ 0) [S(k; A, B)GH(k; \underline{a})]^n S(k; A, B) \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix},\]

where \( S(k; A, B) \) is defined by (3.15). By the unitarity of \( S(k; A, B) \), the series converges absolutely for all \( \underline{a} \in \mathcal{A} \).

Recall that \( S(k; A, B) = S(k; CA, CB) \) for every invertible \( C \). It follows directly from Definition 3.4 that

\[S(k; A, B) = \bigoplus_{v \in V(\mathcal{G})} S(k; A(v), B(v)).\]

Plugging this equality in (4.17) proves the claim. \( \square \)

**Remark 4.11.** Theorem 4.10 implies that the scattering matrix of the graph \( \mathcal{G} \) is determined by the scattering matrices associated with all its single vertex subgraphs. This result can also be obtained by applying the factorization formula [14].

Combining Theorems 4.2 and 4.10 we immediately obtain

**Corollary 4.12.** Let \( \underline{a} \in (\mathbb{R}_+)^{|\mathcal{I}|} \) be arbitrary. For all \( k > 0 \) the scattering matrix \( S(k; A, B, \underline{a}) \) associated with the Laplacian \( \Delta(A, B, \underline{a}) \) on the graph \( \mathcal{G} \) has an absolutely convergent expansion in the form

\[(4.18) \quad S(k; A, B, \underline{a})_{e,e'} = \sum_{w \in \mathcal{W}_{e,e'}} W(w; k) = \sum_{w \in \mathcal{W}_{e,e'}} \hat{W}(w; k)e^{i k |w|}.\]

Since \( \hat{W}(w; k) \) is independent of the metric properties of the graph, it is natural to view (4.18) as a combinatorial Fourier expansion of the scattering matrix \( S(k; A, B, \underline{a}) \).

We will show that (4.18) actually coincides with the Fourier expansion (4.1) in Theorem 4.2.

5. Analytic Continuation of the Scattering Matrix

Recall that the scattering matrix \( S(k; A, B, \underline{a}) \) is analytic in \( k \) for all \( \text{Re} k > 0 \) and \( \text{Im} k > 0 \). In this section we will show that representation (4.18) for the scattering matrix can be extended to the complex plane.
Lemma 5.1. There is $\beta_0 > 0$ such that the series
\[(5.1) \sum_{n \in \mathbb{N}} [\hat{S}_n(k; A, B)]_{e, e'} e^{\overline{n} \cdot \overline{\omega}}\]
converges absolutely for all $k \in \mathbb{C}$ with $\Re k > 0$ and $\Im k \geq \beta_0$. Therefore,
\[
\tilde{S}(k; A, B, \underline{a})_{e, e'} = \sum_{n \in \mathbb{N}} [\hat{S}_n(k; A, B)]_{e, e'} e^{\overline{n} \cdot \overline{\omega}}
\]
is a holomorphic function for all such $k \in \mathbb{C}$.

Proof. From Proposition 3.2 it follows that for all sufficiently large $\beta > 0$ there is a constant $C_\beta > 0$ such that the estimate
\[
|S_v(k)_{e_1, e_2}| \leq C_\beta
\]
holds for all $k \in \mathbb{C}$ with $\Im k > \beta$, any $v \in V$, and any $e_1, e_2 \in E_v$ (see Definition 4.9). Therefore, for an arbitrary walk $w \in W_{e, e'}(\underline{n})$ we obtain
\[
|\tilde{W}(w; k)| \leq C_\beta^{\overline{w} + 1}.
\]
Thus, from (4.12) it follows that
\[
||[\hat{S}_n(k; A, B)]_{e, e'}| \leq C_\beta^{\overline{w} + 1}|W_{e, e'}(\underline{n})|.
\]
Therefore, from Lemma 2.5 using the identity (2.11) we obtain the estimate
\[(5.2) \sum_{n \in \mathbb{N}} |[\hat{S}_n(k; A, B)]_{e, e'}| \leq C_\beta^{N + 1}|\mathcal{I}|^N.
\]
Recalling the definition (2.10) for $a_{\min}$ estimate (5.2) implies that the series
\[
\sum_{n \in \mathbb{N}} |[\hat{S}_n(k; A, B)]_{e, e'}| e^{-n \cdot \Im k a_{\min}} \leq \sum_{N=0}^{\infty} e^{-N \Im k a_{\min}} \sum_{\underline{n} \in \mathbb{N}_{e, e'}} |S(k; \underline{n})_{e, e'}|
\]
converges for all $k \in \mathbb{C}$ with
\[
\Im k > \beta_1 := \frac{1}{a_{\min}} \log \{C_\beta|\mathcal{I}|\}.
\]
This proves the claim with $\beta_0 = \max \{\beta, \beta_1\}$. □

The following statement is the main result of this section.

Theorem 5.2. There is $\beta_0 > 0$ such that
\[(5.3) S(k; A, B, \underline{a})_{e, e'} = \sum_{n \in \mathbb{N}_{e, e'}} [\hat{S}_n(k; A, B)]_{e, e'} e^{\overline{n} \cdot \overline{\omega}}\]
holds for all $k \in \mathbb{C}$ with $\Re k > 0$ and $\Im k > \beta_0$.

The intuitive idea behind the proof of Theorem 5.2 is the observation that the series (4.18) and (5.1) agree. However, Theorem 4.12 and Lemma 5.1 establish convergence of these series in two disjoint sets of the complex plane. Therefore, to prove that both series define the same analytic function we perform a two-step analytic continuation invoking an auxiliary analytic function of two complex variables.
Proof. Consider the $|E| \times |E|$ matrix-valued function $F(k_1, k_2)$ with matrix elements

$$F(k_1, k_2)_{e,e'} := \sum_{n \in N_{e,e'}} [\hat{S}_n(k_1; A, B)]_{e,e'} e^{i k_2 (n,a)}.$$

By Theorem 4.12

$$F(k, k)_{e,e'} = S(k; A, B, a)_{e,e'}$$

for all $k = \text{Re} \, k > 0$ and by Lemma 5.1

$$F(k, k)_{e,e'} = \tilde{S}(k; A, B, a)_{e,e'}$$

for all $k \in \mathbb{C}$ with $\text{Re} \, k > 0$ and $\text{Im} \, k > \beta_0$, where $\beta_0$ is defined in Lemma 5.1. Observe that for any $k_1 > 0$ the function $F(k_1, k_2)$ is holomorphic in $k_2 \in \{k \in \mathbb{C} | \text{Re} \, k > 0, \text{Im} \, k > 0\}$. Assume that $\text{Im} \, k_2 > \beta_0$ with $\beta_0$ defined as in Lemma 5.1. Inspecting the estimates used in the proof of Lemma 5.1 we obtain that

$$\sum_{n \in N_{e,e'}} [\hat{S}_n(k_1; A, B)]_{e,e'} e^{i k_2 (n,a)}$$

converges absolutely for all $k_1 \in \mathbb{C}$ with $\text{Re} \, k_1 > 0$ and $0 \leq \text{Im} \, k_1 < \text{Im} \, k_2 + \epsilon$, where $\epsilon > 0$ is sufficiently small. Recalling (5.4) completes the proof. \qed

Remark 5.3. Assume that the series

$$\sum_{n \in N_{e,e'}} [\hat{S}_n(k; A, B)]_{e,e'} e^{i k_2 (n,a)}$$

absolutely converges in a ball $B_r(k_0)$ centered at $k_0 \in \mathbb{C}$ with $\text{Re} \, k = 0$ and $\text{Im} \, k > 0$. Then arguments used in the proof of Theorem 5.2 show that

$$S(k; A, B, a)_{e,e'} = \sum_{n \in N_{e,e'}} [\hat{S}_n(k; A, B)]_{e,e'} e^{i k_2 (n,a)}$$

for all $k \in B_r(k_0)$. \hfill \qed

6. The Generating Function

In this section we prove an explicit algebraic representation for the matrix-valued generating function $T(\beta)$ defined in equation (2.3). This result is formulated below as Theorem 6.2.

Let $B$ be the canonical orthonormal basis in $\mathbb{C}^{|\mathcal{E}| + 2|\mathcal{I}|} \cong \mathcal{K} = \mathcal{K}_\mathcal{E} \oplus \mathcal{K}_\mathcal{I}^\sim \oplus \mathcal{K}_\mathcal{I}^{+}$ such that any element $h \in B$ is uniquely associated with some edge $j(h) \in \mathcal{I} \cup \mathcal{E}$. Moreover, $j(h) \in \mathcal{E}$ if $h \in \mathcal{K}_\mathcal{E}$ and $j(h) \in \mathcal{I}$ if $h \in \mathcal{K}_\mathcal{I}^\sim$ or $h \in \mathcal{K}_\mathcal{I}^{+}$. Set

$$v(h) = \begin{cases} \partial(j(h)) & \text{if } h \in \mathcal{K}_\mathcal{E}, \\ \partial^-(j(h)) & \text{if } h \in \mathcal{K}_\mathcal{I}^\sim, \\ \partial^+(j(h)) & \text{if } h \in \mathcal{K}_\mathcal{I}^{+}. \end{cases}$$

Given a collection of matrices $\mathcal{M} = \{M(v)\}_{v \in \mathcal{V}}$ we define the linear transformation $M$ on the finite-dimensional Hilbert space $\mathcal{K}$ via its sesquilinear form

$$\langle h_1, Mh_2 \rangle_\mathcal{K} = \begin{cases} [M(v(h))]_{j(h_1),j(h_2)} & \text{if } v(h_1) = v(h_2), \\ 0, & \text{otherwise}. \end{cases}$$
For an arbitrary $\beta > 0$ and every $v \in V(\mathcal{G})$ we set

$$A_v(\beta) := \frac{1}{2}(I - M(v)), \quad B_v(\beta) := -\frac{1}{2\beta}(I + M(v)).$$

Define

$$A(\beta) := \bigoplus_{v \in V} A_v(\beta), \quad B(\beta) := \bigoplus_{v \in V} B_v(\beta).$$

Finally, we set

$$D(\beta) = Z(1\beta; A(\beta), B(\beta), \underline{\underline{a}})$$

$$= \frac{1}{2} \left( X(1\beta; \underline{\underline{a}}) + Y(1\beta; \underline{\underline{a}}) \right) - \frac{1}{2} M \left( X(1\beta; \underline{\underline{a}}) - Y(1\beta; \underline{\underline{a}}) \right)$$

$$= [I + MGH(1\beta; \underline{\underline{a}})]U(1\beta, \underline{\underline{a}}).$$

Here the matrix $Z(k; A, B, \underline{\underline{a}})$ is defined in (3.11), the matrices $X(k; \underline{\underline{a}})$ and $Y(k; \underline{\underline{a}})$ are defined in (3.12), $U(k; \underline{\underline{a}}), G,$ and $H(k; \underline{\underline{a}})$ in (4.15) and (4.16). Writing the matrix $M$ with respect to the orthogonal decomposition (3.3) as a $3 \times 3$ block-matrix

$$M = \begin{pmatrix}
M_{11} & M_{12} & M_{13} \\
M_{21} & M_{22} & M_{23} \\
M_{31} & M_{32} & M_{33}
\end{pmatrix},$$

we obtain

$$I + MGH(1\beta; \underline{\underline{a}}) = \begin{pmatrix}
I & M_{13}e^{-\beta a} & M_{12}e^{-\beta a} \\
0 & I + M_{23}e^{-\beta a} & M_{22}e^{-\beta a} \\
0 & 0 & I + M_{32}e^{-\beta a}
\end{pmatrix}.$$ 

Obviously, $I + MGH(1\beta; \underline{\underline{a}})$ is an entire matrix valued function in the complex variable $\beta$. Moreover,

$$\lim_{\text{Re} \beta \to +\infty} \det(I + MGH(1\beta; \underline{\underline{a}})) = 1.$$ 

Thus, $\det D(\beta)$ is not identically vanishing and this in turn gives

**Lemma 6.1.** The matrix valued function $D(\beta)$ is entire in $\beta \in \mathbb{C}$ and its determinant vanishes on a discrete set $\mathcal{D} \subset \mathbb{C}$ depending on $\underline{\underline{a}} \in \mathbb{R}^{\mathcal{I}}$ and the set of matrices $\mathcal{M} = \{M(v)\}_{v \in V}$. The set $\mathcal{D}$ has no accumulation points in $\mathbb{C}$. In particular, the matrix inverse $D(\beta)^{-1}$ is a meromorphic function in $\beta \in \mathbb{C}$ with poles in $\mathcal{D}$.

Now we turn to the main result of this article:

**Theorem 6.2.** For a given non-compact graph $\mathcal{G} = (V, \mathcal{I}, \mathcal{E}, \vartheta)$ with lengths $\underline{\underline{a}}$ of the internal lines and a collection of matrices $\mathcal{M} = \{M(v)\}_{v \in V}$ at the vertices of the graph the generating function $T(\beta)$ defined by (2.3) has an analytic extension to $\mathbb{C} \setminus \mathcal{D}$ and can be expressed in terms of the matrix $D(\beta)^{-1}M$ as follows

$$T(\beta) = \begin{pmatrix}
I & 0 & 0
\end{pmatrix} D(\beta)^{-1} M \begin{pmatrix}
I \\
0 \\
0
\end{pmatrix}.$$

We turn to the proof of this theorem. First we assume that all matrices $M(v)$ are self-adjoint. Let $A_v(\beta)$ and $B_v(\beta)$ be defined by (6.2). Then $A_v(\beta)B_v(\beta)^{\dagger}$ is self-adjoint, since $M(v)$ is. Observe that

$$\dim \ker (A_v(\beta), B_v(\beta)) = \deg(v) - \dim(\ker A_v(\beta) \cap \ker B_v(\beta)).$$
From (6.2) it follows that \( \text{Ker} A_v(\beta) \cap \text{Ker} B_v(\beta) = \{0\} \). Therefore, the \( 2\deg(v) \times \deg(v) \) matrix \((A_v(\beta), B_v(\beta))\) has maximal rank. Thus, the operator \( \Delta(A_v(\beta), B_v(\beta)) \) for the single-vertex graph \( \hat{G}_v \) (see Definition 4.9) is self-adjoint. The associated scattering matrix given by (3.15) obviously satisfies the relation
\[
S(1\beta; A_v(\beta), B_v(\beta)) = M(v).
\]

Set \( B_r(\beta) = \{ k \in \mathbb{C} | |k - 1\beta| < r \} \).

**Lemma 6.3.** The scattering matrix \( S(k; A_v(\beta), B_v(\beta)) \) is holomorphic for all \( k \in (\mathbb{C}_+ \setminus [0,1\infty)) \cup B_r(\beta) \)

with
\[
r = \frac{\beta}{\|M(v)\|}.
\]

**Proof.** Recalling Proposition 5.2 observe that \( S(k; A_v(\beta), B_v(\beta)) \) has a pole at \( k = 1\pi \), \( \pi \in \mathbb{R}_+ \) if and only if there is a \( \chi \in L_v \) such that
\[
\left( \frac{1}{2} - \frac{\pi}{2\beta} \right) M(v)\chi = \left( \frac{1}{2} + \frac{\pi}{2\beta} \right) \chi,
\]
that is, \((\beta + \pi)(\beta - \pi)^{-1}\) is an eigenvalue of \( M(v) \). Therefore,
\[
\frac{\beta + \pi}{|\beta - \pi|} \leq \|M(v)\|,
\]
which implies that the distance from the point \( 1\beta \) to the closest pole of the scattering matrix \( S(k; A_v(\beta), B_v(\beta)) \) is at least \( \beta \|M(v)\|^{-1} \). \( \square \)

Via equations (3.6) and (3.7) the matrices \( A(\beta), B(\beta) \) being defined by (6.3) define the self-adjoint Laplace operator \( \Delta(A(\beta), B(\beta), a) \) with local boundary conditions (in the sense of Definition 3.4).

Now we choose \( \beta \) so large that the series (2.20) converges. Then, by (6.6), the generating function \( T_{\varepsilon, \varepsilon'}(\beta) \) can represented in the form
\[
T_{\varepsilon, \varepsilon'}(\beta) = \sum_{n \in \mathcal{N}_{\varepsilon, \varepsilon'}} \hat{S}_n(1\beta; A(\beta), B(\beta))_{\varepsilon, \varepsilon'} e^{-\beta(n-a)},
\]
where the coefficients \( \hat{S}_n(1\beta; A(\beta), B(\beta)) \) are defined by (4.11) and (4.12) with the boundary conditions (6.3).

**Lemma 6.4.** Assume that \( \beta > \beta_0 \) with \( \beta_0 \) satisfying (2.12). Then there is \( \rho > 0 \) such that the series
\[
\sum_{n \in \mathcal{N}_{\varepsilon, \varepsilon'}} |\hat{S}_n(k; A(\beta), B(\beta))|_{\varepsilon, \varepsilon'} e^{\beta(n-a)}
\]
converges absolutely for all \( k \in B_{\rho}(\beta) \).

**Proof.** For an arbitrary \( \varepsilon > 0 \) choose \( \rho > 0 \) so small that
\[
|S_v(k)_{e_1, e_2}| \leq \|M(v)\|(1 + \varepsilon)
\]
for all \( k \in B_{\rho}(\beta) \), all \( e_1, e_2 \in \mathcal{E}_v \), and all \( v \in V \). As in the proof of Lemma 5.1 for an arbitrary walk \( w \in \mathcal{W}_{\varepsilon, \varepsilon'}(n) \) we obtain the estimate
\[
|\hat{W}(w; k)| \leq m|w|^{1 + (1 + \varepsilon)|w|^{1}},
\]
where
\[ m := \max_{v \in V} \|M(v)\|. \]

In turn, this implies the bound
\[ \sum_{\mathbf{n} \in N_{e,e'}} |[\hat{S}_{\mathbf{n}}(k; A(\beta), B(\beta))]_{e,e'} e^{i\mathbf{k} \cdot \mathbf{n} \cdot \mathbf{a}}| \leq m^{N+1}(1 + \varepsilon)^{N+1}|I|^N. \]

Therefore,
\[
\sum_{\mathbf{n} \in N_{e,e'}} |[\hat{S}_{\mathbf{n}}(k; A(\beta), B(\beta))]_{e,e'} e^{i\mathbf{k} \cdot \mathbf{n} \cdot \mathbf{a}}| \\
\leq \sum_{N=0}^{\infty} e^{-N \Im k} \sum_{\mathbf{n} \in N_{e,e'}} |[\hat{S}_{\mathbf{n}}(k; A(\beta), B(\beta))]_{e,e'}| \\
\leq \sum_{N=0}^{\infty} e^{-N \Im k} a_{\min} m^{N+1}(1 + \varepsilon)^{N+1}|I|^N.
\]

This series converges if
\[
(6.7) \quad \Im k > \frac{1}{a_{\min}} \log \{m(1 + \varepsilon)|I|\}.
\]

We claim that inequality (6.7) holds for all \( k \in B_\rho(\beta) \) if \( \varepsilon \) is chosen to be so small that
\[
(1 + \varepsilon)e^{\beta_0 - \beta} < 1,
\]
and \( \rho > 0 \) satisfies the inequality
\[
\rho < (\beta - \beta_0) - \frac{1}{a_{\min}} \log(1 + \varepsilon).
\]

Indeed, under these assumptions for any \( k \in B_\rho(\beta) \) we have
\[
\Im k > \beta - \rho > \beta_0 + \frac{1}{a_{\min}} \log(1 + \varepsilon) > \frac{1}{a_{\min}} \log \{m(1 + \varepsilon)|I|\}.
\]

\[ \square \]

\textit{Proof of Theorem 6.2.} Assume the matrices \( M(v) \) to be self-adjoint. Lemma 6.4 and Remark 5.3 imply that there is \( \rho > 0 \) such that
\[
\sum_{\mathbf{n} \in N_{e,e'}} [\hat{S}_{\mathbf{n}}(k; A(\beta), B(\beta))]_{e,e'} e^{i\mathbf{k} \cdot \mathbf{n} \cdot \mathbf{a}} = S(k; A(\beta), B(\beta); \mathbf{a})
\]
holds for all \( k \in B_\rho(\beta) \). Thus, the generating function \( T(\beta) \) can be expressed in terms of the scattering matrix,
\[
(6.8) \quad T(\beta) = S(1\beta; A(\beta), B(\beta); \mathbf{a}).
\]

In turn, the scattering matrix can be calculated by means of Theorem 3.1. Obviously,
\[
Z(1\beta; A(\beta), B(\beta); \mathbf{a}) = D(\beta)
\]
and
\[
A(\beta) + \beta B(\beta) = -M.
\]
Thus, (6.5) follows from (3.14).
Now we relax the assumption on the self-adjointness of the matrices $M(v)$. Obviously, the r.h.s. of (6.3) is a rational function with respect to the entries of the matrix $M$. Since $\det D(\beta)$ does not vanish identically, we obtain the claim. 

\begin{remark}
Relation (6.8) combined with the factorization formula for the scattering matrix on the graph (14) allows to determine the generating function $T(\beta)$ of walks on the graph $G$ in terms of the generating functions associated with subgraphs of $G$.
\end{remark}

\begin{remark}
There is a direct way to establish (6.5). Indeed, observe that by (6.4) one has

\begin{equation}
(I \quad 0 \quad 0) \quad D(\beta)^{-1} M \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix} = (I \quad 0 \quad 0) \quad [I + M \quad G \quad H(1; \underline{a})]^{-1} M \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix}.
\end{equation}

The matrix $G$ performs the “jump” from one boundary vertex of an internal line to the other. A simple calculation shows that if $\text{Re} \quad \beta > 0$, then

$$\|GH(1; \underline{a})\| = e^{-\text{Re} \quad \beta} \quad a_{\text{min}}, \quad \text{where} \quad a_{\text{min}} = \min_{i \in I} a_i.$$

Therefore, the series expansion of $[I + M \quad G \quad H(1; \underline{a})]^{-1}$ converges absolutely for all $\beta \in \mathbb{C}$ with sufficiently large $\text{Re} \quad \beta > 0$. The expression (6.9) coincides with the series (2.3).

This observation gives rise to the following generalization, where the penalty vector depends on the direction in which a given edge is traversed by a walk. Let $\underline{a} = \{a_i\}_{i \in I}$ and $\underline{b} = \{b_i\}_{i \in I}$ be two arbitrary penalty vectors. Set

$$\hat{H}(k; \underline{a}, \underline{b}) = \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix} \quad e^{ik\underline{a}} \quad 0 \\ 0 \\ 0 \quad e^{ik\underline{b}} \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix}.$$

such that $\hat{H}(k; \underline{a}, \underline{a}) = H(k; \underline{a})$. Define now

\begin{equation}
\hat{T}(\beta) = (I \quad 0 \quad 0) \quad [I + M \quad G \quad \hat{H}(1; \underline{a}, \underline{b})]^{-1} M \begin{pmatrix} I \\ 0 \\ 0 \end{pmatrix}.
\end{equation}

For any nontrivial walk $w = \{e', i_1, \ldots, i_N, e\}$ we set

$$c_{ik} = \begin{cases} a_{ik} & \text{if the walk traverses the edge } i_k \in I \text{ in the direction from the terminal to the initial vertex}, \\ b_{ik} & \text{if the walk traverses the edge } i_k \in I \text{ in the direction from the initial to the terminal vertex}, \end{cases}$$

where $k \in \{1, \ldots, \|w\|_{\text{comb}}\}$.

Using the arguments presented above one can easily prove the following statement.

\begin{theorem}
For all $\beta \in \mathbb{C}$ with $\text{Re} \quad \beta$ being sufficiently large the function $\hat{T}(\beta)$ equals the generating function defined by the series (2.3) with $a_{ik}$ being replaced by $c_{ik}$.
\end{theorem}
7. Random Walks on Graphs

In this section we define random walks on a non-compact graph $G$ endowed with the metric structure given by a penalty vector $a$. Assume that the matrices $M(v)$ are stochastic, that is, all their entries are nonnegative and satisfy

$$\sum_{k_1} [M(v)]_{k_1,k_2} = 1 \quad \text{for any edge } k_2 \in I \cup E \text{ incident with the vertex } v,$$

where the sum is taken over all edges $k_1 \in I \cup E$ incident with the vertex $v$. The external lines of the graph will be interpreted as initial or final states of the walk, the internal lines as intermediate states.

Take an arbitrary external line $e \in E$ and consider a sequence $\{X_n\}_{n=0}^N$ of random variables with values in the set $I \cup E$ determined by the following rule. Set $X_0 = e$. Let $v_0 = \partial(e)$. Choose randomly an element $j_1$ of $S(v_0)$ with probability $M(v_0)_{j_1,e}$. Set $X_1 = j_1$. If $j_1 \in E$, then $N = 2$ and the sequence is completed. If $j_1 \in I$, then take $v_1 = \partial(j_1)$, $v_1 \neq v_0$. Choose randomly an element $j_2$ of $S(v_1)$ with probability $M(v_1)_{j_2,j_1}$ and set $X_2 = j_2$. If $j_2 \in E$, then $N = 3$ and the sequence is completed. Otherwise proceed inductively. Finally, we obtain finite or infinite sequence of random variables. If $N < \infty$, then $\{X_n\}_{n=0}^N$ is a walk in the sense of Section 2. We call this sequence a random walk on the graph $G$ from $e \in E$ to $e' = X_N \in E$.

The generating function of random walks from $e \in E$ to $e' \in E$ is defined by equation (2.3). Obviously, it is monotone with respect to $\beta$,

$$T_{e,e'}(\beta) \leq T_{e,e'}(\beta')$$

for $\beta \geq \beta'$. If $W_{e,e'}$ contains at least one nontrivial walk, then $T_{e,e'}(\beta)$ is strictly monotone with respect to $\beta$,

$$T_{e,e'}(\beta) < T_{e,e'}(\beta')$$

for $\beta > \beta'$.

Recall that the stochastic matrix $M(v)$ is said to be regular if it is ergodic, i.e., if there is a natural number $k$ such that the $k$-th power $M(v)^k$ of the matrix $M(v)$ has strictly positive matrix entries.

Lemma 7.1. Let $G$ be a non-compact connected graph. Assume that each $M(v)$ is a regular stochastic matrix. If in addition all diagonal elements of each $M(v)$ are strictly positive, then all matrix elements of $T(\beta)$ are strictly positive for all sufficiently large $\beta > 0$.

Proof. Connectedness of $G$ implies that all $W_{e,e'}$ are non-empty. Given $e$ and $e'$ for $T_{e,e'}(\beta) > 0$ to hold it is necessary and sufficient that there is at least one walk $w \in W_{e,e'}$ with $W(w) > 0$. For the last condition to hold it is in turn sufficient that all matrices $M(v)$ are ergodic and their diagonal elements are strictly positive.

In the remainder of this section we will discuss several examples and introduce some mean values associated with random walks on the graph $G$. These mean values are related to the generating function and its derivative evaluated at $\beta = 0$. However, for $\beta > 0$ the generating function $T_{e,e'}(\beta)$ can be interpreted as a partition function (see, e.g., [23]) with $\beta$ being the inverse temperature. The role of the statistical ensemble is played here by the set $W_{e,e'}(\mathcal{M})$ of all relevant walks from $e'$ to $e$. 

□
1. We leave it to the reader to verify that the mean length of a random walk from \( e' \in \mathcal{E} \) to \( v \in \mathcal{V} \) is given by

\[
\langle |w| \rangle = -\frac{d}{d\beta} \log T_{e,e'}(\beta) \bigg|_{\beta=0} = -T_{e,e'}(\beta)^{-1} \frac{d}{d\beta} T_{e,e'}(\beta) \bigg|_{\beta=0}.
\]

The r.h.s. of (7.1) can be calculated by means of Theorem 6.2. In the thermodynamic setting (i.e., for \( \beta > 0 \)) the quantity

\[ -\frac{d}{d\beta} \log T_{e,e'}(\beta) = -T_{e,e'}(\beta)^{-1} \frac{d}{d\beta} T_{e,e'}(\beta) \]

corresponds to the "mean length" at the temperature \( \beta^{-1} \). In the following examples we will consider probabilistic (\( \beta = 0 \)) and thermodynamic (\( \beta > 0 \)) means on equal ground.

2. As another example we consider the following situation. We say that a walker entering a vertex \( v \) from the edge \( k \) and leaving through the edge \( j \) experiences a transition from \( k \) to \( j \) at \( v \). Now fix a vertex \( v_0 \in \mathcal{V} \) and edges \( j_0, k_0 \in \mathcal{G}_{v_0} \) satisfying the inequality \( M(v_0)_{j_0,k_0} > 0 \). We set

\[
M(v_0; \lambda)_{jk} = \begin{cases} 
  e^{-\lambda} M(v_0)_{jk} & \text{if } j = j_0, k = k_0 \\
  M(v_0)_{jk} & \text{otherwise}
\end{cases}
\]

with an arbitrary \( \lambda > 0 \). Note that \( M(v_0; \lambda) \) is ergodic if \( M(v_0) \) is, but of course not stochastic. Now replacing \( M(v_0) \) by \( M(v_0; \lambda) \) while leaving all other \( M(v) \) in the collection \( \{M(v)\}_{v \in \mathcal{V}} \) unchanged, consider the matrix \( M(\lambda) \) defined by (6.1). Obviously, \( M(0) = M \). Further, similar to (6.4), we introduce the matrix \( D(\beta; \lambda) \)

\[
D(\beta; \lambda) = \frac{1}{2} (X(\beta; a) + Y(\beta; a)) - \frac{1}{2} M(\lambda) (X(\beta; a) - Y(\beta; a)) ,
\]

and define the generating function \( T(\beta; \lambda) \) in analogy with (2.3) by

\[
T_{e,e'}(\beta; \lambda) = \sum_{w \in \mathcal{W}_{e,e'}} W(w; \lambda)e^{-\beta|w|}
\]

with

\[
W(w; \lambda) = \prod_{k=0}^{|w|_{\text{comb}}} [M(v_k; \lambda)]_{e_k^{(+)} e_k^{(-)}} .
\]

Obviously, Theorem 6.2 remains valid for \( T(\beta; \lambda) \) such that

\[
T(\beta; \lambda) = \left( \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right) D(\beta; \lambda)^{-1} M(\lambda) \left( \begin{array}{c} 0 \\ 0 \end{array} \right).
\]

Observe that if \( D(\beta) \) is invertible for a given \( \beta \) then \( D(\beta; \lambda) \) is also invertible for the same \( \beta \) and all sufficiently small \( \lambda > 0 \). We, obviously, have

\[
-\frac{d}{d\lambda} W(w; \lambda) \bigg|_{\lambda=0} = n_{v_0,j_0,k_0}(w) W(w; 0) = n_{v_0,j_0,k_0}(w) W(w),
\]

where \( n_{v_0,j_0,k_0}(w) \geq 0 \) is the number of times a walker experiences a transition from \( k_0 \) to \( j_0 \) at the vertex \( v_0 \) along a given walk \( w \in \mathcal{W}_{e,e'}(\mathcal{M}) \).
Consider the quantity

\[
\langle n_{e,e^\prime}^{v_0,j_0,k_0} \rangle (\beta) = -\frac{d}{d\lambda} \log T_{e,e^\prime}(\beta; \lambda) \bigg|_{\lambda=0}
\]

\[
= -T_{e,e^\prime}(\beta) \left. \frac{d}{d\lambda} T_{e,e^\prime}(\beta; \lambda) \right|_{\lambda=0}.
\]

(7.4)

It is easy to verify that \( \langle n_{e,e^\prime}^{v_0,j_0,k_0} \rangle (\beta) \) is the mean number of times a random walk from \( e^\prime \in E \) to \( e \in E \) experiences a transition from \( k_0 \) to \( j_0 \) at the vertex \( v_0 \).

Using Theorem 6.2 the derivative in (7.4) can be calculated in a rather simple way:

\[
\frac{d}{d\lambda} T(\beta; \lambda) \bigg|_{\lambda=0} = - (I 0 0) \left. \frac{d}{d\lambda} \left( D(\beta; \lambda)^{-1} M(\lambda) \right) \right|_{\lambda=0} \left( \begin{array}{c} I \\
0 \end{array} \right)
\]

\[
= (I 0 0) \left( D(\beta)^{-1} \left. \frac{d}{d\lambda} D(\beta; \lambda) \right|_{\lambda=0} D(\beta)^{-1} M \right) \left( \begin{array}{c} I \\
0 \end{array} \right)
\]

\[
- (I 0 0) \left. \frac{d}{d\lambda} M(\lambda) \right|_{\lambda=0} \left( \begin{array}{c} I \\
0 \end{array} \right).
\]

Now set

\[
\left. -\frac{d}{d\lambda} M(\lambda) \right|_{\lambda=0} = M(v_0, j_0, k_0)
\]

such that

\[
\left. \frac{d}{d\lambda} D(\beta; \lambda) \right|_{\lambda=0} = \frac{1}{2} M(v_0, j_0, k_0) \left( X(1\beta; \underline{a}) - Y(1\beta; \underline{a}) \right),
\]

where the matrices \( X \) and \( Y \) are defined in (3.12). Therefore,

\[
\left. \frac{d}{d\lambda} T(\beta; \lambda) \right|_{\lambda=0} = \frac{1}{2} \left( I 0 0 \right) \left( D(\beta)^{-1} \left( X(1\beta; \underline{a}) - Y(1\beta; \underline{a}) \right) \right)
\]

\[
\cdot M(v_0, j_0, k_0) D(\beta)^{-1} M \left( \begin{array}{c} I \\
0 \end{array} \right)
\]

\[
+ \left( I 0 0 \right) D(\beta)^{-1} M(v_0, j_0, k_0) \left( \begin{array}{c} I \\
0 \end{array} \right).
\]

Thus, only the knowledge of the inverse \( D(\beta)^{-1} \) is necessary to determine \( \left. \frac{d}{d\lambda} T(\beta; \lambda) \right|_{\lambda=0} \).

Note that only one matrix element of \( M(v_0, j_0, k_0) \) is non-vanishing.

The quantities

\[
\langle n_{e,e^\prime}^{v_0,j_0,\bullet} \rangle (\beta) = \sum_{k_0} \langle n_{e,e^\prime}^{v_0,j_0,k_0} \rangle (\beta)
\]

\[
\langle n_{\bullet,j_0,k_0} \rangle (\beta) = \sum_{j_0} \langle n_{v_0,j_0,k_0} \rangle (\beta)
\]
are related to the mean values for the probability that the vertex $v_0$ is entered -- during a walk from $e'$ to $e$ via $j_0 \in S(v_0)$ or left via $k_0 \in S(v_0)$, respectively. Therefore,

$$
\langle n_{e,e'}^{v_0} \rangle (\beta) = \sum_{j_0, k_0 \in S(v_0)} \langle n_{e,e',j_0,k_0}^{v_0} \rangle (\beta)
$$

$$
= \sum_{j_0 \in S(v_0)} \langle n_{e,e',j_0}^{v_0} \rangle (\beta)
$$

$$
= \sum_{k_0 \in S(v_0)} \langle n_{e,e',k_0}^{v_0} \rangle (\beta)
$$

is the mean number of times the vertex $v_0$ is visited during random walks from $e'$ to $e$.

Similarly,

$$
\langle \bar{n}_{e,e'}^{v_0} \rangle (\beta) = \sum_{j_0 \in S(v_0)} \langle n_{e,e',j_0,j_0}^{v_0} \rangle (\beta)
$$

is the mean number of times the vertex $v_0$ is entered and left through the same edge during a walk from $e'$ to $e$.

Assume now that for given $e' \in E$ we have $T_{e,e'} (\beta) > 0$ for all $e \in E$. Set

$$
T_{e,e'} (\beta) = \sum_{e' \in E} T_{e,e'} (\beta).
$$

Then, the value of the quantity

$$
\langle n_{e,e'}^{v_0} \rangle (\beta) = \sum_{e' \in E} \langle n_{e,e'}^{v_0} \rangle (\beta) \frac{T_{e,e'} (\beta)}{T_{e,e'} (\beta)}
$$

gives the mean number of visits at the vertex $v_0$ for random walks starting at $e' \in E$.

Similarly, if for given $e \in E$ we have $T_{e,e'} (\beta) > 0$ for all $e' \in E$ we set

$$
T_{e,e} (\beta) = \sum_{e' \in E} T_{e,e'} (\beta).
$$

The quantity

$$
\langle n_{v_0}^{e,e'} \rangle (\beta) = \sum_{e' \in E} \langle n_{v_0}^{e,e'} \rangle (\beta) \frac{T_{e,e'} (\beta)}{T_{e,e'} (\beta)}
$$

is the mean number of visits of the vertex $v_0$ for walks ending at $e \in E$. With

$$
T_{e,e} (\beta) = \sum_{e' \in E} T_{e,e'} (\beta) = \sum_{e' \in E} T_{e,e'} (\beta) = \sum_{e \in E} T_{e,e} (\beta)
$$

consider the quantity

$$
\langle n_{v_0}^{e,e'} \rangle (\beta) = \sum_{e \in E} \langle n_{v_0}^{e,e'} \rangle (\beta) \frac{T_{e,e} (\beta)}{T_{e,e} (\beta)} = \sum_{e' \in E} \langle n_{v_0}^{e,e'} \rangle (\beta) \frac{T_{e,e'} (\beta)}{T_{e,e'} (\beta)}
$$

$$
= \sum_{e \in E} \langle n_{v_0}^{e,e'} \rangle (\beta).
$$

Obviously, $\langle n_{v_0}^{e,e'} \rangle (\beta)$ is the mean number a random walk in $W(G) = \cup_{e,e'} W_{e,e} (G)$ visits the vertex $v_0$. Therefore,

$$
\sum_{v_0 \in V} \langle n_{v_0}^{e,e'} \rangle (\beta) \geq 1
$$
is the mean number of vertices visited during a random walk.

3. As a final example we consider the mean number \( \langle n_{i_0}^{e'e} \rangle(\beta) \) any internal line \( i_0 \in \mathcal{I} \) is traversed (in either direction) by a random walk from \( e' \) to \( e \). For this replace \( a_{i_0} \) by \( a_{i_0} e^\mu \) while keeping all other \( a_i \) fixed and set
\[
\bar{a}(i_0, \mu) = \{ a_i(i_0, \mu) \}_{i \in \mathcal{I}} \quad \text{with} \quad a_i(i_0, \mu) = \begin{cases} a_i, & \text{if } i \neq i_0, \\ a_{i_0} e^\mu, & \text{if } i = i_0. \end{cases}
\]
Denote by \( T(\beta; \mu) \) the resulting generating function. Then
\[
\langle n_{i_0}^{e'e} \rangle = - \frac{1}{\beta} \frac{d}{d\beta} \log T_{e'e}(\beta; \mu) \bigg|_{\mu=0} = -T_{e'e}(\beta)^{-1} \frac{d}{d\beta} T_{e'e}(\beta; \mu) \bigg|_{\mu=0}.
\]
The derivative of the generating function with respect to \( \mu \) can be calculated by means of Theorem 6.2 thus yielding,
\[
\frac{d}{d\mu} T(\beta; \mu) \bigg|_{\mu=0} = \frac{1}{2} \left[ \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right] \left( D(\beta)^{-1}(I - M) \frac{d}{d\mu} X(1; a(i_0, \mu)) \right) \bigg|_{\mu=0} + (I + M) \frac{d}{d\mu} Y(1; \bar{a}(i_0, \mu)) \bigg|_{\mu=0} \left( D(\beta)^{-1} M \right) \left[ \begin{array}{cc} 1 & 0 \\ 0 & 0 \end{array} \right].
\]

Similar to the discussion of the mean number of vertices visited during a random walk we introduce the quantities
\[
\langle n_{i_0}^{e'e} \rangle(\beta) = \sum_{e' \in E} \langle n_{i_0}^{e'e} \rangle(\beta) \frac{T_{e'e}(\beta)}{T_{e'e}(\beta)},
\]
\[
\langle n_{i_0}^{e'e} \rangle(\beta) = \sum_{e' \in E} \langle n_{i_0}^{e'e} \rangle(\beta) \frac{T_{e'e}(\beta)}{T_{e'e}(\beta)},
\]
\[
\langle n_{i_0}^{e'e} \rangle(\beta) = \sum_{e' \in E} \langle n_{i_0}^{e'e} \rangle(\beta) \frac{T_{e'e}(\beta)}{T_{e'e}(\beta)}.
\]

Thus, \( \langle n_{i_0}^{e'e} \rangle(\beta) \) is the mean number of times the internal line \( i_0 \in \mathcal{I} \) is traversed by a random walk starting at \( e' \in E \), \( \langle n_{i_0}^{e'e} \rangle(\beta) \) the mean number the internal line \( i_0 \in \mathcal{I} \) is traversed by a random walk ending at \( e \in E \). The quantity \( \langle n_{i_0}^{e'e} \rangle(\beta) \) is the mean number of times the internal line \( i_0 \in \mathcal{I} \) is traversed by any random walk.

**APPENDIX. RANDOM WALKS ON VERTICES**

Here we will relate the customary notion of random walks on graphs (see, e.g., [2] or [33]) to random walks considered in the present work. Recall that the customary notion of random walks on graphs is given by a Markov chain with vertices as states. The transition matrix \( P \) indexed by the vertices has a non-vanishing entry only if the corresponding vertices are adjacent.

Consider a graph \( G' = G'(V', \mathcal{I}', \mathcal{E}, \mathcal{O}, \mathcal{B}') \) with no external lines. Let \( P : V' \times V' \rightarrow \mathbb{R}_+ \cup \{0\} \) be a nearest neighbor transition matrix, i.e.,
\[
(A.1) \quad \sum_{v' \in V'} P(v', v) = 1 \quad \text{for any} \quad v \in V
\]
(we read from right to left) and \( P(v', v) > 0 \) occurs only if \( v \) and \( v' \) are adjacent.

Pick an arbitrary vertex in \( G' \) which we denote by \( v_\infty \). Let \( V_{v_\infty} \subset V' \) be the set of all vertices adjacent to \( v_\infty \). Then, \( \I_{v_\infty} \) is the set of the internal lines \( i \in \I \) incident with \( v_\infty \). For any \( i \in \I_{v_\infty} \) let \( v_i \in V_{v_\infty} \) be the vertex adjacent to \( v_\infty \) by \( i \in \I \), that is,

\[
\forall i \in \I_{v_\infty} \quad \partial'(i) = (v_\infty, v_i) \quad \text{or} \quad \partial'(i) = (v_i, v_\infty)
\]

Now replace every edge \( i \in \I_{v_\infty} \) by the external line \( e \) incident with the vertex \( v_i \). Denote the set of all external lines by \( \E \) and define the boundary operator

\[
\partial(j) = \begin{cases} 
\partial'(j), & \text{if } j \in \I', \\
v, & \text{if } j \in \E.
\end{cases}
\]

Thus, we have constructed a non-compact graph \( G(V, \I, \E, \partial) \) with \( V = V' \setminus v_\infty \) and \( \I = \I' \setminus \I_{v_\infty} \). Obviously, the degree of any vertex \( v \in V \) being calculated for the graphs \( G' \) and \( G \) is equal.

Let \( \S(v) \) be the star graph of the vertex \( v \in V \), that is, the set of all edges \( j \in \I \cup \E \) which are incident with the vertex \( v \). Given a matrix \( P \) and \( v \in V \) we define the \( \deg(v) \times \deg(v) \) matrix \( M(v) \) with entries \( M(v)_{ij} \) for \( i, j \in \S(v) \) as follows:

\[
0 \leq M(v)_{ij} = \begin{cases} 
P(v', v) & \text{for } i \in \S(v) \setminus \E, v' \neq v, \\
P(v_\infty, \partial(e)) & \text{for } i = e \in \S(v) \cap \E.
\end{cases}
\]

In particular, the matrix element \( M(v)_{ij} \) is independent of \( j \) and by \( \text{A.1} \)

\[
(A.2) \quad \sum_{i \in \L(v)} M(v)_{ij} = 1 \quad \text{for all } \quad j \in \L(v).
\]

A converse construction is also possible. Assume that a non-compact graph \( G \) has \( \E \neq \emptyset \) and any two vertices of the graph are adjacent by no more than one internal line. Further, assume that all matrix entries \( 0 \leq M(v)_{ij} \) are independent of \( j \), that is, in each matrix \( M(v) \) all columns are equal, and the equality \( \text{A.2} \) holds. Consider the graph \( G' \) without external lines obtained from \( G \) by replacing each external line \( e \) by an internal incident with an additional vertex \( v_\infty \) such that its vertex set \( V' = V \cup \{v_\infty\} \). Now for any \( v', v \in V' \) we set

\[
P(v', v) = \begin{cases} 
M(v)_{ij} & \text{for } i \in \I : v', v \in \partial(i), \\
M(v)_{ij} & \text{for } i \in \E, v' = v_\infty, v = \partial(e), \\
|\E|^{-1} & \text{for } v = v_\infty, v' \in V_{v_\infty}, \\
0 & \text{otherwise}.
\end{cases}
\]

Then \( P \) is a nearest neighbor transfer matrix.

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