Long-term variation in distribution of sunspot groups

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Abstract. We studied the relation between the distribution of sunspot groups and the Gleissberg cycle. As the magnetic field is related to the area of the sunspot groups, we used area-weighted sunspot group data. On the one hand, we confirm the previously reported long-term cyclic behaviour of the sum of the northern and southern sunspot group mean latitudes, although we found a somewhat longer period (\( P \approx 10^4 \) years). We introduced the difference between the ensemble average area of sunspot groups for the two hemispheres, which turns out to show similar behaviour. We also investigated a further aspect of the Gleissberg cycle where while in the 19th century the consecutive Schwabe cycles are sharply separated from each other, one century later the cycles overlap each other more and more.

Key words. Methods: data analysis – Sun: activity – Sun: magnetic fields – sunspots

1. Introduction

The activity of the Sun has been studied in several ways for a long time. The oldest observed and recorded features on the solar disk are the sunspots. Regular sunspot observations were started by Galileo in 1610 and since then the data sets have been being continuously extended. Using these databases one has the opportunity to study the statistical properties of solar activity (see e.g. Vitinsky et al., 1986). The most prominent cycle in the sunspot series and in all solar activity is the 11-year Schwabe cycle. The long-term trend in the amplitude of the Schwabe cycle is known as the secular Gleissberg cycle (Gleissberg, 1945), which has a varying time scale of 80-120 years.

This paper aims to study the relation between the distribution of sunspot groups and the Gleissberg cycle. Although the origin of the Gleissberg cycle is unclear, it is generally accepted that the interaction of the solar differential rotation and the magnetic field plays a basic role in the generation of all solar activity. From the observational data sets a correlation can be seen between secular variations in magnetic activity and differential rotation variations (Vitinsky et al., 1986, Yoshimura & Kambayashi, 1993, Javaraiah, 2003). Although the role of the differential rotation in the long-term variation is not yet clear, several authors have investigated the interplay of differential rotation and the large-scale magnetic field, which drives the solar grand minima. Different modulational mechanisms have been proposed in the recent literature, namely (1) fluctuations in the \( \alpha \)-effect which could cause the dynamo to stochastically flip and lead to modulation of magnetic energy; (2) a dynamic \( \alpha \)-effect forced modulation; (3) two-layer dynamos with meridional circulation; (4) nonlinear dynamics via macroscopic (e.g. Lorentz force) or microscopic (e.g. \( \Lambda \)-quenching) feedback (see a detailed review in Tobias, 2002).

The following study may provide some additional observational background for the investigation of solar dynamo models dealing with the description of the long-term evolution of solar activity.

2. Data and method

We need a detailed and homogenous sunspot catalogue to examine the fine structure of the long-term variation of the sunspot group distribution, so we use the sunspot data of the Greenwich Photoheliographic Results (henceforth GPR) such as the area, the heliographic position and the time of observation.

First, we would like to reanalyse the long term cyclic behaviour of the (signed) sum of the mean latitudes of sunspot groups in the northern and southern hemispheres, respectively, but we weighted the latitudinal distribution of sunspot groups with their area. In a previous work, Pulkkinen et al. (1999) considered 10-day averages of these sums and they found that the “magnetic equator” of the Sun displays a long-term variation between northern and southern sunspot group latitudes with a period of \( P = 33900 \pm 950 \) days, and an amplitude of \( 1.31 \pm 0.13 \) degrees. They obtained similar results with the yearly averages. An important difference with respect to the

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groups. In this case, we do not compose the time average,

time intervals data sets. On the other hand, the average life-
significant differences between the results obtained with dif-
the distributions of sunspot groups deviate from each other. To analyse this discrepancy,
we determine the daily latitudinal area-weighted extension of activity, which can be written in the following form:

\[ E_t = \frac{\sum_{i=1}^{N_t} (\langle A \rangle_t - \Phi_t)^2}{N_t} \]  

Fig. 2 illustrates the daily latitudinal area-weighted extension of activity as a function of time. In this case the discrepancy between the Schwabe minima is more significant, suggesting that as a Gleissberg cycle tends to its maximum, the consecutive cycles overlap each other more and more, which will be discussed in the next section.

3. Results and discussion

3.1. Long-term behaviour of area-weighted mean latitudes of sunspot groups

As it was mentioned in the previous section, we reanalysed the long-term variation of the (signed) sum of the mean latitudes of sunspot groups in the northern and southern hemispheres. An important difference to the work of Pulkkinen et al. (1999) is that we considered the latitudinal distribution of sunspot groups weighted by their area. As a first step of our analysis we applied the discrete Fourier-transform (DFT) method for the data. The main peak belongs to the period of \( P_{DFT} = 42.340 \pm 1.000 \) days. This is significantly larger than the period of Pulkkinen et al. (1999), even larger than the length of our 39 000 day-long observing window. As in such a case the fundamental frequency in the power spectrum is misplaced by aliases, we also searched for the period by phase-dispersion (PHD) method. This resulted in a period of \( P_{PHD} = 37.900 \pm 100 \) days. The difference between the two frequencies is \( |\nu_{DFT} - \nu_{PHD}| = 3 \times 10^{-6} \) day\(^{-1}\). The reason can be understood from the properties of the spectral window of the data. The fundamental frequency of the DFT is \( \nu_{DFT} = 2.3618 \times 10^{-5} \) day\(^{-1}\). The full width at half magnitude of the peak, which depends on the length of the data series, is about \( 2 \times 10^{-5} \) day\(^{-1}\). Consequently, the frequency found by the phase-dispersion method is located very close to the peak. Furthermore, as it is well known, if the spectral window is symmetrical to the origin, a peak presents also at around \(-\nu_{DFT}\). This peak, as well as its side lobes can modify the exact location of the fundamental frequency in the spectrum.

Supposing that the mathematical form of the sinusoidal variation is the following:

\[ f(t) = a_0 + a_1 \cos(2\pi\nu t) + b_1 \sin(2\pi\nu t), \]

we list in Table I the values of the Fourier coefficients obtained by least-squares fits, using the previously found periods, as well as the period given in the paper of Pulkkinen et al. (1999), and the \( \chi^2 \) values obtained by \( \chi^2 \) test.
As we mentioned, the toroidal magnetic flux correlates to the sunspot area, so we construct another dataset in the following way. We take the total area ($A_{30}$) and the number of sunspot groups ($N_{30}$) in every 30 days for the northern and southern hemispheres, respectively. We derived the difference between the ensemble average area of sunspot groups for the two hemispheres, as $(A_{30}/N_{30})_{\text{North}} - (A_{30}/N_{30})_{\text{South}}$. Applying the DFT and the phase dispersion method for these data, the same fundamental period was found (within the error range), namely $P = 48.500 \pm 100$ days. The values of the Fourier-coefficients obtained by least-squares fits are $a_0 = 3.283199(29), a_1 = -6.971889(46)$ and $b_1 = -18.475170(36)$, while the $\chi^2$ value is $\chi^2 = 130.0788751$.

Our results suggest that the toroidal magnetic flux shows a long-term variation between the northern and southern hemispheres, that is, if the ensemble average area of sunspot groups is larger on the southern hemisphere than on the northern one, then we can assume that the toroidal magnetic flux is also larger in the south, and later this dominance alternates.

### 3.3. A further aspect of the Gleissberg cycle

In Figs. [12], it can be seen that the consecutive cycles overlap more and more around the Schwabe minima. In order to study the overlapping areas of the consecutive cycles on a long time scale, we investigated the behaviour of $E_t$ (see Eq. 2) around the minima of 11-year cycles. Thus, we considered two-year intervals centered on the individual minima defined by the relative sunspot number (see review in [Usoskin & Mursula, 2003]). Then we calculated the mean and the standard deviation of the daily latitudinal area-weighted extension of activity ($\langle E \rangle_{\text{min}}$ and $\sigma_{E, \text{min}}$, respectively) for these intervals.

For our study we invoke the results of [García & Mouradian, 1998]. They calculated the secularly smoothed sunspot number of maxima $R'_N(N)$ and minima $R''_N(N)$, where $N$ is the 11-yr solar cycle number and they studied the variation of these computed values in time from the viewpoint of the Gleissberg cycle. Accordingly, their results are well comparable with ours.

In Fig. 5 we present the mean $\langle E \rangle_{\text{min}}$ and the standard deviation ($\sigma_{E, \text{min}}$) of the daily latitudinal area-weighted extension of activity around the minima of 11-yr solar cycle and the results of [García & Mouradian, 1998], namely $R''_N(N)$. It is well visible that as a Gleissberg cycle tends to its maximum $\langle E \rangle_{\text{min}}$ increases. While in the northern hemisphere this tendency is obvious, in the southern one it not so, but the tendency is visible. Both these numbers show that the consecutive cycles overlap each other more and more as the Gleissberg cycle approaches its maximum, i.e. the sunspot groups of the next cycle appear before the end of the previous cycle.

### 4. Conclusion

In this paper we studied the long-term evolution of the distribution of sunspot groups from different aspects, e.g., latitude, area and extension of activity. We applied two new methods, which in our opinion give a more realistic picture of the long-term variation of the toroidal magnetic field. Our method is based on the results of [Petrovay & Szakály, 1999]. They found that...
the latitudinal distribution of the magnetic field at the surface reflects the conditions at the bottom of the convective zone, i.e. in this regard the convective zone behaves as a "steamy window". Although their original statement is restricted to the poloidal magnetic field, it can be assumed that this is also valid for the toroidal component (Petrovay, 2004, private communication). This is why we applied an area-weighted method, since the sunspot area probably reflects the toroidal magnetic flux.

First, we confirmed the results of Pulkkinen et al. (1999) using the area-weighted mean latitude data. However, we found a longer period applying the discrete Fourier-transform and phase-dispersion method, respectively. Based on the \( \chi^2 \) test, the period of \( P_{\text{PHD}} = 37,900 \) days with an amplitude of 1.04° turned out the best. We also studied the difference between the ensemble average area of sunspot groups for the two hemispheres, i.e. \( (A_{30/N30}\text{North} - A_{30/N30}\text{South}) \), which also indicates that the dominance of the toroidal magnetic flux oscillates between the southern and northern hemispheres in a long time scale. For this dataset we found somewhat longer period \( (P = 48,500 \) days). Although the deviation of the two periods is greater than the formal errors (±100 days), the real uncertainty is significantly larger due to the shortness of the observing window with respect to the length of the cycle and the large scatter. Thus, we do not think that the difference of the periods implies real physical effects.

Pulkkinen et al. (1999) interpreted this latitudinal variation as a mixed parity mode in which a quadrupolar component is oscillating with this period. They examined the toroidal components of both the quadrupolar and the dipolar fields based on the paper of Brooke et al. (1998). They claimed that multiplying the absolute value of the total toroidal field by the signed mean latitude of the activity belt, this product should oscillate about zero. As the difference between the ensemble average area of sunspot groups for the two hemispheres, i.e. \( (A_{30/N30}\text{North} - A_{30/N30}\text{South}) \), is related to the strength and pattern of the toroidal magnetic field, our results may indicate observational evidence of such an oscillation. We note that Pipin (1999) also found an asymmetry of magnetic activity.

Beside studying the secular variation of the magnetic activity, we also investigated a further aspect of Gleissberg cycle. In Fig. 10 it is visible that the sequential Schwabe cycles overlap each other at certain times. To characterize the rate of the overlapping we defined the daily latitudinal area-weighted extension of activity. We examined its mean and its standard deviation around the minima of the 11-year solar cycle. These quantities are increasing as the Gleissberg cycle tends to its maximum (see in Fig. 5). While in the northern hemisphere this tendency is obvious, in the southern one it is less so, but the tendency is visible. Nevertheless, as our investigation is limited to one Gleissberg cycle, we cannot declare unambiguously that this phenomenon correlates with the Gleissberg cycle. Pelt et al. (2000) found a similar tendency in the distances in between nearby cycles (see Fig. 11 in Pelt et al., 2000). Hence, the continuous production of detailed, homogenous catalogues (e.g. Debrecen Photoheliographic Data successor of the GPR) is crucial from this point of view. On the other hand, the interpretation of the variable amount of overlap between the Schwabe cycle is a challenge for future solar dynamo models.

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Fig. 5. The mean (solid line) and the standard deviation (dashed line) of the daily latitudinal extension of activity as a function of time. The dotted line represents the results of Garcia & Mouradian (1998), i.e. the secularly smoothed sunspot number of minima \( R_m \).
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