Deposition of Crystalline GdIG Samples Using Metal Organic Decomposition Method

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Abstract: Fabrication of high quality ferrimagnetic insulators is an essential step for ultrafast magnonics, which utilizes antiferromagnetic exchange of the ferrimagnetic materials. In this work, we deposit high-quality GdIG thin films on a (111)-oriented GGG substrate using the Metal Organic Decomposition (MOD) method, a simple and high throughput method for depositing thin film materials. We postannealed samples at various temperatures and examined the effect on structural properties such as crystallinity and surface morphology. We found a transition in the growth mode that radically changes the morphology of the film as a function of annealing temperature and obtained an optimal annealing temperature for a uniform thin film with high crystallinity. Optimized GdIG has a high potential for spin wave applications with a low damping parameter in the order of $10^{-3}$, which persists down to cryogenic temperatures.

Keywords: garnet ferrite; compensated ferrimagnet; metal organic decomposition

1. Introduction

High frequency dynamics from the antiferromagnetic exchange are essential for achieving fast computing spintronic devices [1,2]. On the other hand, to utilize the antiferromagnetic dynamics, a method to overcome its lack of responsiveness to the external field needs to be devised. Compensated ferrimagnet, composed of two antiferromagnetically coupled sublattices [3], enables the fast dynamics of antiferromagnetic materials while having accessibility via external fields similar to ferromagnetic ones.

The spin wave carries information as a collective precessional motion of localized magnetization. Unlike electric current mediated by the motion of charge carriers, the spin wave does not involve joule heating when propagating. Thus, it has high potential in the applications such as low-power computing devices [3]. However, a low damping constant is required to ensure a long transport distance for practical application. Rare earth iron garnet (REIG) is one of the best candidates in ferrimagnetic magnon transport. As garnet ferrite is an insulator, there is no electron spin sink contribution to damping, which results in low damping [4].

One of the obstacles in using REIG is a complicated fabrication process [5], which requires a dedicated vacuum chamber [6] or annealing temperature near the melting point [6–8]. The metal organic decomposition (MOD) method provides a simple and reproducible method to deposit REIG. Unlike other methods, the MOD method only needs
a spin coating of the inorganic compound solution and the annealing temperature lower than the melting point.

In this work, we prepared gadolinium iron garnet (GdIG) thin film on gadolinium gallium garnet (GGG) substrate using the MOD method. X-ray diffraction (XRD), transmission electron microscopy (TEM), superconducting quantum interference device vibrational sample magnetometer (SQUID VSM), and Atomic force microscope (AFM) were applied to check crystallinity, magnetic properties, and morphology of samples. We found that high-quality garnet samples with reproducible magnetic properties can be deposited, which can be used to study physical phenomena such as longitudinal spin Seebeck effect [9] and spin Hall effect-driven various magnetoresistances [10,11].

2. Materials and Methods

The precursor solution is prepared by dissolving inorganic compounds Gd(NO$_3$)$_3$·6H$_2$O with 99.99% purity (from Alfa Aesar, Ward Hill, MA, USA) and Fe(NO$_3$)$_3$·6H$_2$O, which has 99.9% purity (from Sigma Aldrich, St. Louis, MO, USA), into the solvent. A mixture of 99.8% purity dimethylformamide (DMF/Sigma Aldrich, St. Louis, MO, USA) and 95% purity polyvinylpyrrolidone (PVP) was applied as a solvent [12]. DMF and PVP were chosen because the mixture can dissolve inorganic compounds and, hence, can be utilized to grow garnet samples using MOD method [12–14]. As the mixing rate of DMF and PVP can dominantly affect the thickness and morphology of the sample, we used the reported values for YIG and BiYIG [13] because they have similar structures with GdIG. The stoichiometry of the Fe compound and Gd compound was set to 3:5, based on the formula unit of Gd$_3$Fe$_5$O$_{12}$. The total concentration of metallic compounds was fixed to 17%.

In order to avoid the contamination at surface of substrate, the GGG (111) substrate was sonicated first in acetone and then in ethanol for 30 min each. After that, the sample was treated under Ar plasma (20 sccm of Ar, 100 @ power with 70 kHz) for 30 min. A MOD solution was spin-coated onto GGG (111) substrate at 500 rpm for 5 s and 3000 rpm for 30 s. The coating and solution preparing procedure was optimized to deposit 10 nm of garnet sample [12]. For thicker samples, we repeated the growth procedure multiple times. After coating, the sample was dried for 30 min at 100 °C using a hot plate. The sample was then annealed for 1 h under an oxygen-rich atmosphere at various temperatures (750 °C, 800 °C, 900 °C, and 1000 °C) to determine the optimal annealing temperature for high crystallinity and good surface morphology. A fixed ramping rate of 4 °C/min was used to reach the target temperature.

Crystallinity was confirmed using a high-resolution XRD (RIGAKU corporation) with K$_\alpha$ emission line of Copper (λ = 1.54 Å). Surface morphology was examined by using AFM (XE-7 designed by Park Systems). Magnetic property, especially the magnetization compensation temperature ($T_M$), was measured by SQUID VSM in the Quantum Design magnetic property measurement system (MPMS).

3. Results

3.1. Structural Properties

To find the optimal post-annealing temperature for crystal formation, we first deposited a GdIG sample, targeting 10 nm thickness on the GGG substrate. Figure 1a,b show the measured XRD spectrum of deposited GdIG samples for various annealing temperatures. GGG (444) peak appears at 51.2°, which indicates that the lattice constant of GGG is 12.37 Å. Next to GGG (444) peak, a broad peak is observed at 50.3°, which corresponds to the GdIG (444) peak with a lattice constant of 12.55 Å. We note that the lattice constant of GdIG on GGG is larger than that of a single crystalline GdIG of 12.48 Å due to the pseudo-morphic growth on the GGG substrate [15]. We confirmed that there is no extra peak in the entire XRD pattern (Figure 1a) except for 111 orientation, suggesting the high epitaxial quality of GdIG. The thickness of GdIG was confirmed using a fringe pattern in the XRD spectrum. For the 750 °C annealed sample, the first and second fringe peaks
were observed at 52.08° and 52.98°, from which the thickness of GdIG was estimated to be 11 ± 2 nm. This confirms that the thickness of GdIG is almost the same as what we targeted.

Figure 1. XRD spectrum for GdIG samples annealed at various annealing temperatures (a) from 20° to 90° and (b) from 48.5° to 53°. XRD spectrum for GdIG samples annealed at various annealing times (c) from 20° to 90° and (d) from 48.5° to 53°.

As shown in Figure 1b, the increased annealing temperature drives GdIG peak shaper. As the broadening of XRD peak is caused by vacancy- or defect-driven lattice expansion, the XRD result suggests that the higher annealing temperature results in better crystallinity [5]. We also confirmed that annealing time does not exert any significant change in the XRD pattern (Figure 1c). This suggests that crystallinity is predominantly governed by the annealing temperature.

However, we find that the high annealing temperature can affect the surface topography. The surface topography map measured using AFM in the non-contact mode (Figure 2) shows a significant change in surface morphology as annealing temperature increases. The roughness of film surface is about 0.1 nm order for the 750 °C and 800 °C annealed samples (see Table 1). However, the surface morphology of GdIG changes abruptly above 900 °C, and the island starts to appear. For 1000 °C annealed samples, island growth becomes the dominant film formation mechanism, resulting in the entire area being covered with islands of ~200 nm diameter and ~20 nm height.

Table 1. Roughness of films (Ra) for several annealing temperatures.

| Annealing Temperature (°C) | Ra (nm) |
|---------------------------|--------|
| 750                       | 0.14   |
| 800                       | 0.08   |
| 900                       | 0.71   |
| 1000                      | 8.71   |

As the purpose of this work is to deposit uniform film for spin wave application, we selected 750 °C postannealed sample and performed additional measurement using TEM to ensure the crystallinity. TEM measurement was conducted using the STEM (scanning transmission electron microscope) system with Cs-corrector (JEM-2100F) designed by JEOL. A high-resolution TEM image in Figure 3a shows that GdIG and GGG have a smooth interface with a continuous crystal structure above and below the interface due to a similar lattice constant between the two materials. Additionally, a clear separation of Fe in the
GdIG film and Ga in the GGG substrate from energy dispersive X-ray spectroscopy (EDX) (Figure 3b) indicates minimal mixing between GdIG and substrate.

Figure 2. AFM image of (a) 750 °C, (b) 800 °C, (c) 900 °C, and (d) 1000 °C annealed samples.

Figure 3. (a) TEM and (b) EDX images of 750 °C annealed GGG | GdIG sample. As GGG and GdIG have similar structure and lattice parameters, a smooth interface appears, indicating an epitaxial growth of GdIG. (b) The presence of Gd and Fe and the absence of Ga at the marked region represent the formation of the GdIG layer.

3.2. Magnetic Property

The magnetic property was checked as an additional measure to ensure crystallinity and stoichiometry. Temperature dependent magnetometry of 10 nm GdIG sample annealed at 750 °C was performed using SQUID VSM in MPMS 3. An M-T curve was obtained by measuring M-H curves under in-plane magnetic field at various temperatures and extracting the saturation magnetization ($M_s$). We eliminated paramagnetic contribution from GGG substrate by subtracting the linear background in the M-H curve at each temperature.
Figure 4 shows the extracted M-T curve. We focus on the value of $M_S$ at low temperatures and the magnetic compensation point. $M_S$ at 20 K was about 700 emu/cc, similar to previous reports as shown in Table 2. Magnetic compensation point appeared near 270 K also agrees with references on bulk and thin film GdIG [9,16,17].

![Figure 4. M-T curve of GdIG 10 nm thin film post-annealed at 750 °C. Inset shows the M-H curve at 20 K.](image)

Table 2. Compensation temperature $T_M$ and saturation magnetization ($M_s$) at 20 K of GdIG.

| Source          | $T_M$  | $M_s$ (emu/cc) |
|-----------------|--------|---------------|
| From Figure 4   | 270 K  | 700           |
| From [16]       | 280 K  | 600           |
| From [17]       | 300 K  | 670           |
| From [9]        | 280 K  | 1300          |

3.3. FMR Measurement

Finally, the damping parameter of the GdIG sample was measured using the temperature dependent FMR. We used 60-nanometer-thick GdIG sample to enhance RF absorption power. Temperature and field control was performed using a physical property measurement system (PPMS). RF field for FMR was applied using a custom-built FMR setup that is compatible with PPMS.

Figure 5a shows the FMR spectra of the GdIG film for various temperatures. Here, we fixed the excitation frequency (15 GHz) and swept the magnetic field from 0 to 5000 Oe along the in-plane direction. Figure 5b shows resonance frequency as a function of external magnetic field for $T = 50$ K. The exact overlap with Kittel formula (red solid line) was observed. We extracted the damping parameter from the spectrum by measuring peak to peak distance from the FMR lineshape in Figure 5a. We first use the model for conventional ferromagnetic material to obtain damping parameter $\alpha_{FM}$:

$$\Delta H = \Delta H_0 + \frac{4\pi a_{FM}}{\sqrt{3} \gamma_{eff}} f$$

where $\gamma_{eff} = g_{eff} \mu_B / h = \Delta M / \Delta S$ is the effective gyromagnetic ratio of the net moment, and $\Delta M$ and $\Delta S$ are the differences in the net moment and the spin density of two sublattices, respectively. Figure 5c shows $\Delta H$ as a function of frequency for several temperatures. The slope increases with increasing temperature. Considering that the temperature variation of $\gamma_{eff}$ is small, as seen in Table 3, the result suggests that $a_{FM}$ increased with increasing temperature, which we ascribed to the divergence of $a_{FM}$ at angular momentum compensation point ($T_A$) [18–20].
Figure 5. (a) FMR spectra of GdIG 60 nm samples at 15 GHz RF frequency. The horizontal axis (field axis) is shifted as much as the resonance field to compare the linewidth for different temperatures. (b) Resonance frequency as a function of the external field at 50 K and the Kittel formula fitting (red line). (c) Frequency dependence of peak-to-peak FMR linewidth. The solid line represents linear fitting. (d) Temperature dependence of damping parameters $\alpha_{FM}$ and $\alpha_{FiM}$.

Table 3. Effective gyromagnetic ratio ($\gamma_{eff}$) and damping parameter ($\alpha_{FM}$) measured at various temperatures below $T_M$.

| Temperature (K) | $\gamma_{eff}$ ($\times 10^7$ T$^{-1}$ s$^{-1}$) | $\alpha_{FM}$ ($\times 10^{-3}$) |
|----------------|-----------------------------------------------|-------------------------------|
| 50             | 1.78 ± 0.08                                   | 1.60 ± 0.22                   |
| 100            | 1.85 ± 0.07                                   | 2.80 ± 0.22                   |
| 150            | 1.81 ± 0.02                                   | 4.67 ± 0.74                   |
| 200            | 1.75 ± 0.08                                   | 7.41 ± 0.3                    |

In the case of ferrimagnet, however, the dissipation rate from the two lattices and the dependence of resonance frequency on net spin density need to be considered to deduce a well-defined damping constant [18–21]. Since the GdIG sample also can be regarded as a compensated ferrimagnet with two effective sublattices, damping $\alpha_{ FiM}$ is expressed as follows:

$$\alpha_{ FiM} = \frac{|s_{net}|}{|s_{total}|} \alpha_{ FM}$$

where $s_{net} = s_1 - s_2$ and $s_{total} = s_1 + s_2$ are the difference and the sum of the spin density for each sublattice ($s_{1,2}$), respectively.

To estimate $\alpha_{ FiM}$, we should extract the spin density of each sublattice in GdIG. This was achieved by using the M-T curve shown in Figure 4 and by assuming that the spin density of the Fe sublattice follows the same trend with that in YIG [16]. As shown in Figure 5d, $\alpha_{ FiM}$ had no singular behavior close to angular momentum compensation in contrast to $\alpha_{ FM}$, and this dependence corresponds to previous reports on $\alpha_{ FM}$ of compensated ferrimagnets, including metallic GdFeCo [20].

We note that the GdIG sample deposited by the MOD method has comparable damping constant to some of the literature on spin wave applications [22,23]. In the case of the 10 nm YIG sample grown by the Pulsed Laser Deposition (PLD) method, an order of $10^{-3}$ damping parameter was obtained, which is similar to our result [24]. Our data also
suggest that the MOD-deposited GdIG can be an alternative platform that could be used for low-temperature magnonic applications.

4. Conclusions

A high-quality GdIG thin film on GGG substrate was prepared using the MOD method. We found that a transition of the growing model of GdIG occurs at around 900 °C above which the film grows with an island pattern. The optimal postannealing temperature was decided based on the growth mode of the film and the crystallinity confirmed by the TEM. Temperature dependence of magnetization was measured using SQUID VSM and confirmed that $T_M$ was measured at 270 K. FMR measurement showed that MOD deposited GdIG exhibits comparable damping constant to various deposition methods. Moreover, our experiment found that the damping constant of GdIG decreased at lower temperatures. We expect that the MOD method provides a simple and high-throughput procedure to deposit GdIG, which has potential advantages for magnonic applications.

Author Contributions: K.-J.K. and J.-R.J. planned and supervised the study. P.-C.V. deposited the GdIG sample; H.J. and J.-W.Y. performed the FMR measurement. J.Y. measured AFM; Y.J. measured MPMS; H.K. and A.M.P analyzed the data from AFM, XRD, FMR, TEM, EDX, and MPMS and wrote the manuscript. All authors were involved in the discussion of the results and commented on the manuscript. All authors have read and agreed to the published version of the manuscript.

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Abbreviations

The following abbreviations are used in this manuscript:

- MOD Metal Organic Decomposition;
- FMR Ferromagnetic resonance;
- REIG Rare Earth Iron garnet;
- YIG Yttrium Iron garnet;
- GdIG Gadolinium Iron garnet;
- GGG Gadolinium Gallium Garnet;
- AFM Atomic Force Microscopy;
- XRD X-ray Diffraction;
- TEM Transmission of Electron Microscope;
- EDX Energy Dispersive X-ray spectroscopy;
- PVP PolyVinylPyrrolidone;
- DMF DiMethylFormamide;
- SQUID Superconducting Quantum Interference Vibrating System;
- VSM Vibrational Sample Magnetometer;
- PPMS Physical Property Measurement System;
- MPMS Magnetic Property Measurement System;
- PLD Pulsed Laser Deposition.


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