Proton induced dielectron radiation off Nb: $P_l$ and $Y$ distributions

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Abstract. Following our recent letter [1] on inclusive $e^+e^-$ pair production in proton induced reactions at $E_{\text{kin}} = 3.5$ GeV on the nucleus Nb, we present here in addition the transverse and rapidity distributions for various $e^+e^-$ invariant mass bins and compare them to reference data measured in p+p reactions.

1. Introduction
The prospect to relate hadron properties inside a strongly interacting medium with chiral symmetry restoration has motivated plenty of work, both in theoretical and experimental physics. The spontaneous breaking of chiral symmetry leads to a non vanishing value of the two-quark condensate, which is related to the dynamical generation of hadron masses. Focussing mainly on vector mesons, various models predict relatively strong changes of particle masses and/or widths already at normal nuclear matter density $\rho_0$ [2, 3, 4, 5]. However a consistent picture of in-medium hadron properties has not yet emerged and more experimental input is needed. Moreover such modifications are expected to be most pronounced for particles with small relative momenta to the surrounding medium according to hadronic models [6], a region which is challenging to access in experiment.

Experimentally, in-medium properties can be directly studied using the decay of vector mesons into an $e^+e^-$ pair, since electrons and positrons are not affected by strong final state interactions. Although medium modifications are expected to be stronger in heavy-ion collisions, compared proton-, pion- or photon- induced reactions on nuclei, measured observables represent an average over the complete space-time evolution of the temperature and the density of the system and are hence more complicated to model. On the other hand, in induced reactions on nuclei the system does not undergo a noticeable density and temperature evolution in time, hence conditions of the system are better defined. The experimental drawback is that, for a measurement sensitive to the in-medium spectral shape, the decay to an $e^+e^-$ pair has to take place inside the nucleus. Therefore good acceptance for decays of low momentum vector mesons is of crucial importance, in particular, for the relatively long living $\omega$ and $\phi$ mesons.

Figure 1. Dielectron cross sections measured in p+Nb and p+p reactions at a kinetic beam energy of 3.5 GeV. In case of the p+p data a comparison to a calculated PYTHIA/PLUTO dielectron cocktail is shown in addition (the different sources are listed in the legend).
In fact, most experiments focusing on the spectral distribution of dielectrons produced off nuclei in photon and proton induced reactions are restricted to relatively high momenta ($P_{ee} > 0.8$ GeV/c) and are not conclusive yet. For the $\rho$ meson, the CLAS experiment at JLab [7] reports a slight broadening and no shift of the $\rho$ pole position in photon induced reactions, while the E325 experiment at KEK [8] deduced a shift but no broadening in proton induced reactions. In this paper we follow our recent letter [1] on inclusive e$^+$e$^-$ pair production in proton induced reactions at $E_{\text{kin}} = 3.5$ GeV on the nucleus Nb, representing the first high statistics measurement with small momenta of e$^+$e$^-$ pairs relative to the medium ($P_{ee} < 0.8$ GeV/c). We present here in addition the transverse momentum and rapidity distributions for various e$^+$e$^-$ invariant mass bins. These data are compared to our reference data measured in p+p reactions at the same incident beam energy [9].

2. Experimental data

The efficiency corrected invariant mass distributions of e$^+$e$^-$ pairs within the HADES acceptance are displayed in Fig. 1. The systematic uncertainties are represented by colored horizontal, which are the quadratic sum of errors estimated from different particle identification methods (10%), from consistency checks of the efficiency correction (10%), including the uncertainty due to combinatorial background subtraction, and from the normalization to cross sections (15%). The total systematic error amounts to 21% in case of the p+Nb data, and to 20% for the p+p data. Only the systematic errors of the normalization are taken into account for the comparison of the shape of the invariant mass spectra as the other systematic errors should cancel at first order. For the p+p data, a dielectron cocktail, generated using an adapted version of the event generator PYTHIA (see [9] for details) is displayed in addition. We distinguish between four different mass regions: $M_{ee}$ [GeV/c$^2$] $< 0.15$ (dominated by neutral pion decays), $0.15 < M_{ee}$ [GeV/c$^2$] $< 0.47$ ($\eta$ Dalitz decay dominated), $0.47 < M_{ee}$ [GeV/c$^2$] $< 0.7$ (dominated by direct $\rho$ decays and Dalitz decays of baryonic resonances and $\omega$ mesons) and $0.7 < M_{ee}$ [GeV/c$^2$] ($\rho$ and $\omega$ dominated) as can be seen from the cocktail. Around 1 GeV/c$^2$ a low statistics $\phi$ signal is observed.

1 Note that, according to the systematics on dielectron emission obtained in p+p and d+p collisions at various beam energies by the DLS collaboration [10], it is safe to conclude that at the kinetic beam energy of 3.5 GeV isospin effects play only a secondary role. Therefore, in order to extract medium effects in p+Nb, the above discussed p+p data at the same kinetic beam energy of 3.5 GeV/c is a valuable reference.
Figure 3. Transverse momentum distributions of dielectrons in p+p (scaled) and p+Nb collisions.

visible, which will be further discussed in future, making use of additional information from its hadronic decay channel to charged kaons. The underestimation of the dielectron yield in the mass region from $0.47 < M_{ee} [GeV/c^2] < 0.7$ is addressed in e.g. [12, 13], where the authors improve the description by the introduction of a coupling between $\rho$ mesons and baryonic resonances.

Comparing the shape of the invariant mass spectra separately for pairs with momenta larger and smaller than 0.8 GeV/c to pairs from p+p, scaled to the number of participants and the total reaction cross section, we observe a strong $e^+e^-$ excess yield below the $\omega$ pole mass in the small momentum sample displayed in the right panel of Fig. 2, while for pairs with $P_{ee} > 0.8$ GeV/c no significant difference in the vector meson mass region within the systematic uncertainties indicated by the colored bands in the panel of Fig. 2, is visible. For pairs with $P_{ee} < 0.8$ GeV/c (right panel of Fig. 2) the $e^+e^-$ yield at the $\omega$ pole mass is not reduced, but as the underlying smooth distribution is enhanced, the yield in the peak is reduced to almost zero within errors. This additional yield we attribute to $\rho$-like channels as it is supposed to be the dominating source for radiation from the medium due to its large total width. For more details, see [1].

Comparing the $P_t$ and $Y$ distributions for the four previously discussed invariant mass regions the very same pattern is observed: In the p+Nb system the distributions reach out to slightly higher transverse momentum and the rapidity distributions shift towards target rapidity, see Fig. 3 and Fig. 4. The latter one is much stronger pronounced and corresponds to second generation particles produced from a slower source compared to the initial one. Although this behavior is to some extent expected (stronger transformation of longitudinal energy to transverse degrees of freedom in p+Nb) the defined values are important constraints to phenomenological models in order to correctly model the experimental conditions.
3. Summary

In addition to our recent letter [1] on inclusive $e^+e^−$ pair production in proton induced reactions at $E_{kin} = 3.5$ GeV on the nucleus Nb, we presented the transverse and rapidity distributions for various $e^+e^−$ invariant mass bins. Comparing these distributions to reference data from p+p collisions, the following pattern is observed: In the p+Nb system the distributions reach out to slightly higher transverse momentum and the rapidity distributions are shifted towards target rapidity. Especially the defined shift in rapidity gives important constraints to phenomenological models.

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