Molecular gas mass functions of normal star-forming galaxies since $z \sim 3$*

S. Berta$^1$, D. Lutz$^1$, R. Nordon$^2$, R. Genzel$^1$, B. Magnelli$^3$, P. Popesso$^1$, D. Rosario$^1$, A. Saintonge$^1$, S. Wuyts$^3$, and L. J. Tacconi$^1$

1 Max-Planck-Institut für extraterrestrische Physik (MPE), Postfach 1312, 85741 Garching, Germany
e-mail: berta@mpe.mpg.de
2 School of Physics and Astronomy, The Raymond and Beverly Sackler Faculty of Exact Sciences, Tel-Aviv University, 69978 Tel-Aviv, Israel
3 Argelander-Institut für Astronomie, Universität Bonn, Auf dem Hügel 71, 53121 Bonn, Germany

Received 26 April 2013 / Accepted 7 June 2013

ABSTRACT

We used deep far-infrared data from the PEP/GOODS-Herschel surveys and restframe ultraviolet photometry to study the evolution of the molecular gas mass function of normal star-forming galaxies. Computing the molecular gas mass, $M_{\text{mol}}$, by scaling star formation rates through depletion timescales, or combining infrared (IR) luminosity and obscuration properties as described in the literature, we obtained $M_{\text{mol}}$ for roughly 700, $z = 0.2$–3.0 galaxies near the star-forming main sequence. The number density of galaxies follows a Schechter function of $M_{\text{mol}}$. The characteristic mass $M^*$ is found to strongly evolve up to $z \sim 1$ and then flatten at earlier epochs, resembling the IR luminosity evolution of similar objects. At $z \sim 1$, our result is supported by an estimate based on the stellar mass function of star-forming galaxies and gas fraction scalings from the PHIBSS survey. We compared our measurements with results from current models, finding better agreement with those that are treating star formation laws directly rather than in post-processing.

Integrating the mass function, we studied the evolution of the $M_{\text{mol}}$ density and its density parameter $\Omega_{\text{mol}}$.

Key words. galaxies: luminosity function, mass function – galaxies: statistics – galaxies: evolution – galaxies: star formation – infrared: galaxies

1. Introduction

Stars form in cold, dense molecular clouds, and the molecular gas content of galaxies is an important constraint to galaxy evolution. In the interplay between accretion of gas, star formation, metal enrichment, and outflows, the molecular gas content reflects the prevailing physical processes (e.g. Bouché et al. 2010; Davé et al. 2010; Lilly et al. 2013). For the cosmic baryon budget (e.g. Fukugita et al. 1998), molecular gas can be of increasing relative importance at high redshifts where galaxies are gas-rich.

Molecular gas is generally quantified using the CO molecule as a tracer. Local samples include several hundred targets (e.g. Saintonge et al. 2011, and references therein), but CO detections of star-forming, intermediate- or high-redshift galaxies are still limited to modest statistics of mostly very luminous galaxies (Carilli & Walter 2013, and references therein). For normal star-forming galaxies, the PHIBSS survey (Tacconi et al. 2013) derives scaling relations on the basis of CO detections of 52 normal star-forming galaxies at $z \sim 1.2$ and $z \sim 2.2$. In galaxies near the star-forming main sequence (MS, e.g. Noeske et al. 2007; Elbaz et al. 2007; Daddi et al. 2007), the molecular gas mass, $M_{\text{mol}}$, scales as star formation rate (SFR) through a depletion timescale $\tau_{\text{dep}}$ that only very weakly depends on redshift (Tacconi et al. 2013).

Keres et al. (2003) reported the first attempt of deriving the local molecular gas mass function, based on a sample of IRAS-selected objects. To date, no derivation of a $z > 0$ mass function has been possible due to the heterogeneity and paucity of the available detections.

These difficulties have sparked interest in methods that use dust mass as a tool for deriving masses of the cold ISM of (distant) galaxies (e.g. Magdis et al. 2011, 2012; Scoville 2012). Dust mass is converted into gas mass by adopting a metallicity and scaling the gas-to-dust ratio with metallicity. In order to derive accurate dust masses, these methods require photometry on the restframe-submm tail of the spectral energy distributions (SED) or accurate multi-band data close to the restframe far-infrared (FIR) SED peak. These requirements are still not met individually for large samples of normal high-$z$ galaxies (Berta et al., in prep.). In a study of the restframe ultraviolet (UV) and FIR properties of Herschel galaxies, Nordon et al. (2013) established a method for deriving $M_{\text{mol}}$ of MS galaxies on the basis of their rest-FIR luminosity and rest-UV obscuration: these properties are readily available for larger samples. Here we apply the $\tau_{\text{dep}}$ scaling and the Nordon et al. recipe to the deepest Herschel (Pilbratt et al. 2010) FIR extragalactic observations and related restframe UV data to derive $M_{\text{mol}}$ of MS galaxies, which are known to power $\sim 90\%$ of the cosmic star formation (Rodighiero et al. 2011). On this basis, we construct their molecular gas mass function and study its evolution since $z \sim 3$.

We adopt a $\Lambda$CDM cosmology with $(h,\Omega_m,\Omega_\Lambda) = (0.70, 0.27, 0.73)$ and a Chabrier (2003) initial mass function.

2. Derivation of molecular gas mass and source selection

We used the deepest PACS (Poglitsch et al. 2010) FIR survey, combining the PACS Evolutionary Probe (PEP, Lutz et al. 2011)
realizations of $M_{\text{mol}}$ for each galaxy assuming Gaussian distributions for the mentioned sources of errors and accounting also for the contribution of the template library intrinsic scatter to the $L$(IR) uncertainty. As a result, the median uncertainty on $M_{\text{mol}}$ obtained through Eq. (1) is ~40%, without accounting for the Tacconi et al. (2013) systematics, and exceeds 50% for only ~10% of the sample.

Sample selection is driven by the need to compute total SFRs and by the requirements imposed by Eq. (1). We limited the analysis to galaxies lying within $|\Delta \text{log} (\text{SFR})_{\text{obs}}| \leq 0.5$ from the MS. The MS was assumed to have unit slope in the stellar mass vs. star formation rate ($M_*$ − SFR) plane, and a specific-SFR normalization varying as $sSFR_{\text{MS}}$ (Gyr$^{-1}$) = $26 \times T_2^{0.2}$ (Elbaz et al. 2011).

We applied a 3σ flux cut at 160 μm, the band that best correlates with $L$ (IR) (Elbaz et al. 2011; Nordon et al. 2012). When deriving UV parameters, it is important to avoid contamination of observed bands by the 2100 Å carbonaceous absorption feature. Combining GALEX and optical photometry, we defined four redshift windows: $0.2 < z < 0.6$, $0.7 < z < 1.0$, $1.0 < z < 2.0$, and $2.0 < z < 3.0$. This choice reduces the loss of sources due to restframe UV requirements to <2%. The total number of sources in each redshift bin is included in Table 1. Our sample contains 43 galaxies hosting an X-ray active galactic nucleus (AGN), of which only four are Type-1. The distributions of $L$ (IR) and $\beta$ for the AGN galaxies are similar to those of inactive ones. We verified that our conclusions below are not affected by the inclusion of these AGN hosts.

3. Molecular gas mass function

The comoving number density of galaxies in intervals of $M_{\text{mol}}$ was computed adopting the well-known $1/V_s$ formalism:

$$\Phi(M) \Delta M = \sum \frac{1}{V_s} \Delta M_i,$$  \hspace{1cm} (2)

where the sum is computed over all sources in the given $M_{\text{mol}}$ bin. The accessible volume $V_s$ is a spherical shell delimited by $z_{\text{bin}}$ and $z_{\text{min}}$, $z_{\text{max}}$. Here $z_{\text{max}}$ is the highest redshift at which a galaxy would be observable in our survey (Schmidt 1968), and $z_{\text{min}}, z_{\text{max}}$ define redshift bins.

Our source selection was mainly based on a 160 μm cut, and the UV requirements do not produce significant source losses. On the other hand, the conversion from $S$ (160) to $L$ (IR) was obtained by adopting a family of SED templates spanning a variety of colors. Moreover, UV properties are also involved in computing the total SFR and in applying Eq. (1).

As a consequence, the conversion between the 160 μm fluxes and molecular gas mass is not unique, but comprises a distribution of mass-to-light ratios. Because of this scatter in $M_{\text{mol}}$ vs. $S$ (160), a flux cut induces a molecular mass incompleteness. This effect was thoroughly studied by Fontana et al. (2004) and Berta et al. (2007), among others, for stellar mass. Equivalently, the recipe defined by these authors can be applied to the specific case of molecular gas mass and FIR fluxes to derive completeness corrections as a function of $M_{\text{mol}}$. Using the distribution of the parent MS population and comparing it with that of PACS galaxies leads to similar results.

The comoving number density of galaxies is shown in Fig. 1. Results from the $t_{\text{dep}}$ scaling or Eq. (1) are very similar, thus only the latter are shown. Table 1 reports the two estimates, along with average completeness values for each mass bin, which also includes the photometric completeness of the FIR catalogs.
Table 1. Molecular gas mass function of Herschel galaxies derived with the 1/$V_v$ method.

| log $(M_{mol})$ | $0.2 < z \leq 0.6$ | $0.7 < z \leq 1.0$ | $1.0 < z \leq 2.0$ | $2.0 < z \leq 3.0$ |
|----------------|------------------|------------------|------------------|------------------|
|                | $\Phi(M)$ | $\Phi(M)$ | Compl. | $\Phi(M)$ | $\Phi(M)$ | Compl. | $\Phi(M)$ | $\Phi(M)$ | Compl. |
| $[10^9 h_{70}^2 M_\odot]$ | $\tau_{av}$ | $[10^{-4} h_{70}^2$ Mpc$^{-3}$ dex$^{-1}$] | $\tau_{av}$ | $[10^{-4} h_{70}^2$ Mpc$^{-3}$ dex$^{-1}$] | $\tau_{av}$ | $[10^{-4} h_{70}^2$ Mpc$^{-3}$ dex$^{-1}$] | $\tau_{av}$ | $[10^{-4} h_{70}^2$ Mpc$^{-3}$ dex$^{-1}$] |
| 9.0            | 23.42 ± 6.26 | 0.52 | – | – | – | – | – | – | – |
| 9.2            | 25.38 ± 4.80 | 18.42 ± 5.11 | 0.67 | – | – | – | – | – | – |
| 9.4            | 22.01 ± 4.02 | 24.42 ± 4.39 | 0.78 | – | – | – | – | – | – |
| 9.6            | 13.00 ± 2.77 | 17.86 ± 3.32 | 0.86 | – | – | – | – | – | – |
| 9.8            | 8.06 ± 2.08 | 16.85 ± 2.98 | 0.98 | 12.84 ± 3.87 | 20.15 ± 6.08 | 0.61 | – | – | – |
| 10.0           | 6.97 ± 1.86 | 10.46 ± 2.28 | 1.00 | 10.67 ± 1.98 | 13.53 ± 2.02 | 0.85 | 10.47 ± 2.18 | 15.47 ± 2.82 | 0.57 |
| 10.2           | 2.49 ± 1.11 | 4.48 ± 1.49 | 1.00 | 11.43 ± 1.79 | 10.52 ± 1.64 | 0.96 | 7.05 ± 1.09 | 7.27 ± 1.01 | 0.72 |
| 10.4           | 1.99 ± 1.00 | 1.99 ± 1.00 | 1.00 | 8.33 ± 1.43 | 6.37 ± 1.25 | 1.00 | 7.09 ± 0.85 | 5.37 ± 0.65 | 0.82 |
| 10.6           | 1.00 ± 0.70 | 0.50 ± 0.50 | 1.00 | 3.74 ± 0.97 | 4.17 ± 1.01 | 1.00 | 3.60 ± 0.47 | 2.65 ± 0.38 | 0.90 |
| 10.8           | 0.50 ± 0.50 | – | 1.00 | 1.32 ± 0.59 | 1.22 ± 0.55 | 1.00 | 1.67 ± 0.30 | 1.64 ± 0.28 | 1.00 |
| 11.0           | – | – | – | 0.98 ± 0.49 | – | 1.00 | 0.67 ± 0.18 | 0.38 ± 0.13 | 1.00 |
| 11.2           | – | – | – | – | – | – | 0.35 ± 0.13 | 0.28 ± 0.12 | 1.00 |
| 11.4           | – | – | – | – | – | – | 0.05 ± 0.05 | 0.05 ± 0.05 | 1.00 |
| 11.6           | – | – | – | – | – | – | 0.05 ± 0.05 | 0.05 ± 0.05 | 1.00 |
| 11.8           | – | – | – | – | – | – | 0.04 ± 0.04 | 0.08 ± 0.06 | 0.00 |

Tot. Num. 145 166 260 122

Fig. 1. Molecular gas mass function of Herschel galaxies. Left: comparison of the 1/$V_v$ estimate (green squares, based on Eq. (1)) with literature data (Keres et al. 2003) and models (Obreschkow & Rawlings 2009a; Lagos et al. 2011; Duffy et al. 2012). The red line and shaded area at $z = 1.0$–2.0 are obtained by scaling the Muzzin et al. (2013) stellar mass function using the molecular gas fractions reported by Tacconi et al. (2013). When needed, masses found in the literature were scaled by the factor 1.36 necessary to account for helium, and were matched to our set of cosmological parameters. Right: results of the parametric STY evaluation of a Schechter mass function (red lines) and its 3σ uncertainty (shaded areas).

An independent characterization of the mass function is provided by the maximum-likelihood approach by Sandage et al. (1979, STY). We adopted the Bayesian implementation by Berta et al. (2007) and adapted it to our case, thus fully propagating the actual $M_{mol}$ uncertainties into the parametric function evaluation. The adopted functional form to describe the mass function is a Schechter (1976) function:

$$
\Phi(M)\, dM = \frac{\Phi^*}{M^\alpha} \left(\frac{M}{M^*}\right)^{-\alpha} \exp\left(-\frac{M}{M^*}\right)\, dM,
$$

where $\Phi^*$ represents the normalization, $\alpha$ the slope in the low-mass regime, and $M^*$ the transition mass between a power-law and the exponential drop-off, and the e-folding mass of the latter. Berta et al. (2007) provided more details about this method. Table 2 includes the most probable parameter values and their 3σ uncertainties, obtained with both $M_{mol}$ estimates. The right-hand panel of Fig. 1 compares the result of the STY analysis with the 1/$V_v$ approach. No STY analysis was attempted for the highest redshift bin because only the very massive end is covered.

The Schechter $M^*$ parameter increases by more than a factor of 3 between $z = 0.4$ and 0.8, and then flattens at $z > 1$. This rate resembles the evolution of the IR luminosity function of normal star-forming galaxies (Gruppioni et al. 2013) and reflects the link between SFR and $M_{mol}$. At the same time, $\Phi^*$ varies by only a factor of 2 over the 0.4–1.5 redshift range. The net effect...
Table 2. Results of the STY analysis and molecular gas mass density. Top: results based on Eq. (1). Bottom: results obtained with the $\tau_{dep}$ scaling of SFR.

| 0.2 < z ≤ 0.6 | 0.7 < z ≤ 1.0 | 1.0 < z ≤ 2.0 | 2.0 < z ≤ 3.0 |
|----------------|---------------|---------------|---------------|
| Based on Eq. (1) | Based on Eq. (1) | Based on Eq. (1) | Based on Eq. (1) |
| Most prob. | Min. | Max. | Most prob. | Min. | Max. | Most prob. | Min. | Max. | Most prob. | Min. | Max. |
| $\Phi^{\ast}$ [10$^{-2}$ $h_0^2$ Mpc$^{-3}$ dex$^{-1}$] | 11.0 | 9.8 | 11.0 | 7.8 | 7.4 | 10.6 | 5.6 | 5.3 | 6.6 | – | – | – |
| $\log(M^\ast)$ [10$^{10}$ $M_\odot$] | 13.34 | 10.23 | 10.34 | 10.84 | 10.72 | 10.86 | 10.89 | 10.84 | 10.94 | – | – | – |
| $\alpha$ | –0.91 | –1.08 | –0.90 | –1.12 | –1.15 | –1.03 | –1.16 | –1.16 | –1.04 | – | – | – |
| $\rho_{\rm mol}$(00) [10$^3$ $h_0^2$ M$_{\odot}$ Mpc$^{-3}$] | 2.266 | 1.751 | 2.289 | 5.974 | 4.317 | 7.818 | 4.836 | 4.112 | 5.711 | – | – | – |
| $\rho_{\rm mol}$(lim) [10$^3$ $h_0^2$ M$_{\odot}$ Mpc$^{-3}$] | 2.166 | 1.619 | 2.188 | 4.173 | 3.231 | 5.361 | 3.958 | 3.301 | 4.972 | 1.533 | 1.184 | 1.882 |
| $\Omega_{\rm mol}$(00) [10$^{-7}$ $h_0^2$] | 11.34 | 8.76 | 11.46 | 18.00 | 13.01 | 23.56 | 7.19 | 6.12 | 8.50 | – | – | – |
| $\Omega_{\rm mol}$(lim) [10$^{-7}$ $h_0^2$] | 10.84 | 8.10 | 10.95 | 12.57 | 9.74 | 16.15 | 5.89 | 4.91 | 7.40 | 0.92 | 0.71 | 1.13 |

| 0.2 < z ≤ 0.6 | 0.7 < z ≤ 1.0 | 1.0 < z ≤ 2.0 | 2.0 < z ≤ 3.0 |
|----------------|---------------|---------------|---------------|
| Based on $\tau_{dep}$ | Based on $\tau_{dep}$ | Based on $\tau_{dep}$ | Based on $\tau_{dep}$ |
| Most prob. | Min. | Max. | Most prob. | Min. | Max. | Most prob. | Min. | Max. | Most prob. | Min. | Max. |
| $\Phi^{\ast}$ [10$^{-2}$ $h_0^2$ Mpc$^{-3}$ dex$^{-1}$] | 12.0 | 9.8 | 13.6 | 7.4 | 6.8 | 9.3 | 4.9 | 4.3 | 6.2 | – | – | – |
| $\log(M^\ast)$ [10$^{10}$ $M_\odot$] | 10.32 | 10.27 | 10.36 | 10.79 | 10.70 | 10.83 | 10.87 | 10.83 | 10.93 | – | – | – |
| $\alpha$ | –1.01 | –1.11 | –1.00 | –1.06 | –1.10 | –1.00 | –1.11 | –1.14 | –1.04 | – | – | – |
| $\rho_{\rm mol}$(00) [10$^3$ $h_0^2$ M$_{\odot}$ Mpc$^{-3}$] | 2.540 | 1.963 | 3.114 | 4.725 | 3.640 | 6.287 | 3.880 | 3.205 | 5.407 | – | – | – |
| $\rho_{\rm mol}$(lim) [10$^3$ $h_0^2$ M$_{\odot}$ Mpc$^{-3}$] | 2.400 | 1.810 | 2.943 | 3.437 | 2.771 | 4.407 | 3.232 | 2.590 | 4.698 | 1.213 | 0.929 | 1.497 |
| $\Omega_{\rm mol}$(00) [10$^{-7}$ $h_0^2$] | 12.71 | 9.82 | 15.59 | 14.21 | 10.97 | 18.94 | 5.77 | 4.77 | 8.04 | – | – | – |
| $\Omega_{\rm mol}$(lim) [10$^{-7}$ $h_0^2$] | 12.01 | 9.06 | 14.73 | 10.36 | 8.35 | 13.28 | 4.81 | 3.85 | 6.99 | 0.73 | 0.56 | 0.90 |

Notes. Minimum and maximum values are computed at the 3σ confidence level. The molecular gas mass density $\rho$(tot) is obtained by integrating the mass function between 10$^7$ $M_\odot$ and infinity, $\rho$(lim) is computed solely on the mass range covered by PEP data (see Table 1).

is a significant evolution of the number density of galaxies with large $M_{\rm mol}$, while at the low-mass end it remains roughly constant. Finally, the most probable value of $\alpha$ steepens as redshift increases, but this might be simply an effect of the different mass ranges effectively constrained at different redshifts (note that the 3σ confidence levels are consistent with nearly no evolution).

We compared this $M_{\rm mol}$ mass function with an estimate based on stellar mass. Using the PHIBSS CO survey, Tacconi et al. (2013, see their Fig. 12) computed molecular gas fractions for $z = 1.0$–1.5 normal star-forming galaxies as a function of $M_\ast$. We combined these gas fractions with the $z = 1.0$–1.5 stellar mass function of star-forming galaxies by Muzzin et al. (2013, see also Lib et al. 2013; Drory et al. 2009). To account for the MS width in this estimate, we applied a 0.2 dex Gaussian smoothing. The result is close to the observed mass function at $z = 1.0$–2.0 (Fig. 1). Note also the related approach of Sargent et al. (in prep.), as quoted in Carilli & Walter (2013).

4. Discussion

Our observed $M_{\rm mol}$ function was computed only for normal star-forming galaxies within ±0.5 dex in SFR from the main sequence. Thus it represents a lower limit to the total $M_{\rm mol}$ mass function, by missing passive (low SFR) galaxies and powerful, above-sequence objects. Passive galaxies will provide a negligible contribution unless they include a hypothetical population of molecular gas-rich galaxies forming stars with very low efficiency. For reference, local passive galaxies, lying at $\Delta$(SFR)$_{MS}$ ≤ –1.0 dex, have molecular gas fractions $f_{\rm mol} \leq 2\%$ (Saintonge et al. 2012). Star-bursting galaxies, i.e., those lying above the MS, should play only a minor role in shaping the mass function as well Rodighiero et al. (2011) have shown that the contribution of above-sequence sources to the number density of star-forming galaxies and total SFR density at $z = 1.5$–2.5 is low. For our $\Delta$log (SFR)$_{MS} = 0.5$ cut, we found that objects above the MS contribute no more than ~4% to the number density of star-forming galaxies and ~16% to their SFR density. Depending on whether galaxies rise above the main sequence mostly due to an increased star formation efficiency at fixed gas mass or due to larger gas masses, the effect on the mass function will vary between a global upward ~4% shift or a preferential increase at higher gas masses within the limits permitted by the 16% SFR contribution. Saintonge et al. (2012) showed that in local starbursts the two effects share a 50%–50% role in causing SFR changes with respect to the MS.

Our results are compared in Fig. 1 with the Keres et al. (2003) local CO mass function of IRAS-selected galaxies. Applying a variable CO-to-H$_2$ conversion factor, inversely dependent on L(CO) itself, Obreschkow & Rawlings (2009b,a) obtained a revisited mass function that quickly drops at the high-mass end. Modeling was implemented by post-processing the De Lucia & Blaizot (2007) semi-analytic (SAM) results and assigning the atomic and molecular gas content to galaxies via a set of physical prescriptions (see Obreschkow et al. 2009). Their expectations (Fig. 1) tend to predict a too steep mass function at both low and high redshift. A second model based on the SAM approach was developed by Lagos et al. (2011), starting from the Bower et al. (2006) galaxy formation model. As in Obreschkow & Rawlings (2009a), the Blitz & Rosolowsky (2006) star formation law was adopted, but in this case it was implemented throughout the galaxy evolution process in the SAM model. The model was tested against the observed stellar mass density evolution and the atomic and molecular gas content of local galaxies. Results are shown in Fig. 1, and are now much closer to our observed mass function up to $z \sim 2$. Figure 1 finally includes predictions of the hydrodynamical simulation developed by Duffy et al. (2012), which overall tend to overestimate the mass function, but are consistent with data at the observed high-mass tail.

Adopting the STY Schechter results, we integrated the mass function from 10$^7$ $M_\odot$ to infinity and obtained a measure of the molecular gas mass density (see Table 2 and top panel of Fig. 2). A second estimate obtained by limiting the integral to the mass range effectively covered by observations is also provided. The molecular gas mass density increases by a factor
of $\sim$4 from $z = 0$ to $z = 1$ and then remains almost constant up to $z \sim 2$. For reference we plot also the $M_\star$ density (Muzzin et al. 2013; Ilbert et al. 2013) up to $z = 4$. The different trends reflect the growth of gas fractions as a function of redshift (e.g., Tacconi et al. 2013).

Finally, we computed the redshift evolution of the density parameter $\Omega_{\text{mol}}(z) = \rho_{\text{mol}}(z)/\rho_c(z)$ (Fig. 2 and Table 2), where the critical density at the given redshift is given by $\rho_c(z) = 3H(z)^2/8\pi G$. The molecular gas density parameter peaks at $z \sim 1.0$; for comparison $\Omega_H$, derived from damped Ly-$\alpha$ systems (Prochaska & Herbert-Fort 2004; Rao et al. 2006), and including also low column density cases, does not evolve between $z = 0.5$ and 4.0.

Using the deepest Herschel extragalactic observations available and restframe UV information for roughly 700 main-sequence galaxies, we have built the first molecular gas mass function at redshift $z > 0$ and the first estimate of its density evolution up to $z = 3$. While future mm/submm surveys will significantly improve our knowledge of the molecular content in high-$z$ galaxies, we have provided a basis for refinement of galaxy evolution models that are accounting for the molecular phase.

Acknowledgements. PACS has been developed by a consortium of institutes led by MPE (Germany) and including UVIE (Austria); KU Leuven, CSL, IMEC (Belgium); CEA, LAM (France); MPIA (Germany); INAF-IASF/OAA/OAP/OAT, LENS, SISSA (Italy); IAC (Spain). This development has been supported by the funding agencies BMVIT (Austria), ESA-PRODEX (Belgium), CEA/CNES (France), DLR (Germany), ASI/INAF (Italy), and CICYT/MCYT (Spain).

References

Berta, S., Lonsdale, C. J., Polletta, M., et al. 2007, A&A, 476, 151
Berta, S., Magnelli, B., Nordon, R., et al. 2011, A&A, 532, A49
Berta, S., Lutz, D., Santini, P., et al. 2013, A&A, 551, A100
Blitz, L., & Rosolowsky, E. 2006, ApJ, 650, 933
Bouché, N., Dekel, A., Genzel, R., et al. 2010, ApJ, 718, 1001
Bower, R. G., Benson, A. J., Malbon, R., et al. 2006, MNRAS, 370, 645
Carilli, C., & Walter, F. 2013 [arXiv:1301.6371]
Chabrier, G. 2003, PASP, 115, 763
Daddi, E., Dickinson, M., Morrison, G., et al. 2007, ApJ, 670, 156
Dave, R., Finlator, K., Oppenheimer, B. D., et al. 2010, MNRAS, 404, 1355
De Lucia, G., & Blaizot, J. 2007, MNRAS, 375, 2
Dror, N., Bundy, K., Leauthaud, A., et al. 2009, ApJ, 707, 1595
Duffy, A. R., Kay, S. T., Battye, R. A., et al. 2012, MNRAS, 420, 2799
Elbaz, D., Daddi, E., El Borgne, D., et al. 2007, A&A, 468, 33
Elbaz, D., Dickinson, M., Hwang, H. S., et al. 2011, A&A, 533, A119
Fontana, A., Pozzetti, L., Donnarumma, I., et al. 2004, A&A, 424, 23
Fukugita, M., Hogan, C. J., & Peebles, P. J. E. 1998, ApJ, 503, 518
Grazian, A., Fontana, A., de Santis, C., et al. 2006, A&A, 449, 951
Gruppioni, C., Pozzi, F., Rodighiero, G., et al. 2013, MNRAS, 432, 23
Ilbert, O., McCracken, H. J., Le Fevre, O., et al. 2013, A&A, in press, 10.1051/0004-6361/201321100
Kennicutt, Jr., R. C. 1998, ARA&A, 36, 189
Keres, D., Yun, M. S., & Young, J. S. 2003, ApJ, 582, 659
Lagos, C. D. P., Baugh, C. M., Lacey, C. G, et al. 2011, MNRAS, 418, 1649
Lilly, S. J., Carollo, C. M., Pippino, A., Renzini, A., & Peng, Y. 2013 [arXiv:1303.5959]
Lutz, D., Poglitsch, A., Aliieri, B., et al. 2011, A&A, 532, A90
Magdis, G. E., Daddi, E., Elbaz, D., et al. 2011, ApJ, 740, L15
Magdis, G. E., Daddi, E., Bethermin, M., et al. 2012, ApJ, 760, 6
Magnelli, B., Elbaz, D., Chary, R. R., et al. 2011, A&A, 528, A35
Magnelli, B., Poppingo, P., Berta, S., et al. 2013, A&A, 553, A132
Meurer, G. R., Heckman, T. M., & Calzetti, D. 1999, ApJ, 521, 64
Muzzin, A., Marchesini, D., & Santini, P., et al. 2013, ApJ, 745, 182
Nordon, R., Lutz, D., Genzel, R., et al. 2012, ApJ, 745, 182
Nordon, R., Lutz, D., Saintonge, A., et al. 2013, ApJ, 762, 125
Obreschkow, D., & Rawlings, S. 2009a, ApJ, 696, L129
Obreschkow, D., & Rawlings, S. 2009b, MNRAS, 394, 1857
Obreschkow, D., Croton, D., De Lucia, G., Khochfar, S., & Rawlings, S. 2009, ApJ, 698, 1467
Pilbratt, G. L., Riedinger, J. R., Passvogel, T., et al. 2010, A&A, 518, L1
Poglitsch, A., Waiaekens, C., Geis, N., et al. 2010, A&A, 518, L2
Prochaska, J. X., & Herbert-Fort, S. 2004, PASP, 116, 622
Rao, S. M., Turnshek, D. A., & Nestor, D. B. 2006, ApJ, 636, 610
Rodighiero, G., Daddi, E., Barollo, E., et al. 2011, ApJ, 739, L40
Saintonge, A., Kauffmann, G., Kramer, C., et al. 2011, MNRAS, 415, 32
Saintonge, A., Tacconi, L. J., Fabello, S., et al. 2012, ApJ, 758, 73
Sandage, A., Tammann, G. A., & Yahil, A. 1979, ApJ, 232, 352
Santini, P., Fontana, A., Grazian, A., et al. 2009, A&A, 504, 751
Schechter, P. 1976, ApJ, 203, 297
Schmidt, M. 1968, ApJ, 151, 393
Scoville, N. Z. 2012 [arXiv:1219.6999]
Tacconi, L. J., Neri, R., Genzel, R., et al. 2013, ApJ, 768, 74
Teplitz, H. I., Chary, R., Elbaz, D., et al. 2011, AJ, 141, 1
Wyithe, S., Förster Schreiber, N. M., Lutz, D., et al. 2011, ApJ, 738, 106
Zwaan, M. A., Meyer, M. J., Staveley-Smith, L., & Webster, R. L. 2005, MNRAS, 359, L30

Fig. 2. Top: redshift evolution of the molecular gas mass density based on Eq. (1). Red symbols and solid error bars belong to the total mass density, dashed error bars mark lower limits limited to the mass range covered by observations. The black star is computed by integrating the local mass function (Keres et al. 2003). Gray squares and triangles represent the total $M_\star$ density reported by Ilbert et al. (2013) and Muzzin et al. (2013), respectively. Light-blue symbols belong to star-forming galaxies only. Bottom: evolution of $\Omega_{\text{mol}}$ (red symbols, this work) and $\Omega_H$ (Prochaska & Herbert-Fort 2004; Zwaan et al. 2005; Rao et al. 2006), where we have divided the latter by 1.3 to avoid counting helium twice.