PSR B1951+32: A BOW SHOCK–CONFINED X-RAY NEBULA, A SYNCHROTRON KNOT, AND AN OPTICAL COUNTERPART CANDIDATE

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ABSTRACT

The radio pulsar PSR B1951+32 and the supernova remnant CTB 80 provide a rich laboratory for the study of neutron stars and supernova remnants. Here, we present ground-based optical and near-infrared observations of them, along with X-ray observations with Chandra and a reanalysis of archival data obtained with the Hubble Space Telescope. The X-ray observations reveal a cometary pulsar wind nebula that appears to be confined by a bow shock produced by the high-velocity motion of the pulsar, making PSR B1951+32 a rare pulsar exhibiting both an Hα bow shock and a shocked X-ray pulsar wind nebula. The distribution of the nebula and radio continuum emission is indicative of a contact discontinuity of the shocked pulsar winds and shocked ambient medium at ∼0.05 pc. On the other hand, the optical synchrotron knot of PSR B1951+32 likely has a flat spectrum in the optical and near-infrared wave bands, and our astrometry is consistent with only one of the two reported optical counterpart candidates for the pulsar.

Subject headings: pulsars: individual (PSR B1951+32) — shock waves — stars: neutron — supernova remnants

1. INTRODUCTION

Rotation-powered pulsars (RPPs) exhibit diverse interesting phenomena ranging from periodic pulsations to supersonic motions through the interstellar medium as well as relativistic winds often revealed by interaction with the neighborhood. More than 30 years after the discovery of RPPs, however, the origins for many of these phenomena still remain out of our grasp, although there has been a recent surge of development in some of these fields, largely driven by new X-ray satellites typified by Chandra. One of the many unprecedented features of the new X-ray satellites is superb angular resolutions; for example, Chandra gives ≈1″ accuracy in the positions of X-ray sources, comparable to the typical optical and near-infrared (IR) observations. This enables astronomers to make high-precision comparative studies of X-ray and optical/near-IR results, such as more robust searches for optical/near-IR counterparts to RPPs and multiwavelength studies of their close vicinity.

In many ways, the ∼39.5 ms radio pulsar PSR B1951+32 (Kulkarni et al. 1988) in the CTB 80 supernova remnant (SNR) provides a rare opportunity to study the interesting phenomena of RPPs. It is a moderately young pulsar with a spin-down age of 1.1 × 10^5 yr, which is comparable to, but slightly larger than, the age determined by its proper motion (6.4 × 10^4 yr; Migliazzato et al. 2002) and the dynamical age of CTB 80 (7.7 × 10^4 yr; Koo et al. 1990), indicating a finite initial spin. The measured proper motion (25 ± 4 mas yr^{-1}) corresponds to a velocity of 240 ± 40 km s^{-1} at 2 kpc, moving away from the SNR center. In the optical, PSR B1951+32 is located in an ∼1′ nebular core emitting both Balmer-dominated and forbidden lines. Hester & Kulkarni (1988, 1989) suggested that the core is a bow shock formed by the relativistic winds from the pulsar when it encountered the material behind a radiative shock of the SNR. ROSAT X-ray observations were also suggestive of the pulsar wind nature of the core, albeit its inadequate spatial resolution (Safi-Harb et al. 1995). On the other hand, using archival data from the Hubble Space Telescope (HST), Hester (2000) reported a possible optical synchrotron knot near PSR B1951+32, and Butler et al. (2002) claimed the detection of its two optical counterpart candidates. However, the lack of color information and the rather large astrometric uncertainties make it difficult to reach any further conclusion on their nature.

In this Letter, we present X-ray and optical/near-IR observations of PSR B1951+32, searching for the predicted X-ray pulsar wind nebula in the optical core as well as the emission from the claimed optical synchrotron knot and counterpart candidates.

2. OBSERVATIONS AND RESULTS

The CTB 80 core was observed with Chandra on 2001 July 12 for a total good time of 78.2 ks using its Advanced CCD Imaging Spectrometer (S3 chip), with a custom subarray mode of 192 rows, minimizing pileup on the pulsar while still imaging the entire core. Figure 1 shows a resulting X-ray image, where a strong point source (i.e., PSR B1951+32; ∼21,900 photon counts) is embedded in diffuse emission. Its position is R.A. = 19^h52^m58^s, decl. = +32°52′41″ with 0.02 uncertainties11 in each coordinate (2000). Considering the 0.7′ systematic uncertainty of the Chandra astrometry, the X-ray position is consistent with that inferred from pulsar radio timing (Foster et al. 1994). Additionally, several other point sources were detected in the Chandra image, and

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11 All positional uncertainties quoted in this Letter represent the 90% confidence levels.
Fig. 1.—Chandra X-ray (0.3–10 keV) image (both in gray scale and contour) of the CTB 80 core. The image was convolved with a Gaussian filter of 1° FWHM. The contour levels correspond to 0.1%, 0.2%, 0.4%, 0.8%, 1.1%, 1.5%, 1.9%, 2.6%, 3.8%, 5.7%, 10.5%, and 19% of the peak brightness.

Fig. 2.—Chandra X-ray image (contour) superposed on the Hα image obtained with the HST F656N filter (gray scale). The axes represent the offsets from the pulsar position. The contour levels are 0.1%, 0.4%, 0.7%, 1.3%, 2.2%, 4.0%, 7.3%, 36.6%, and 95.2% of the peak brightness for the upper panel and 0.9%, 1.2%, 1.6%, 2.2%, 3.1%, 4.4%, 7.3%, 22.0%, and 73.2% for the bottom panel. The HST image was convolved with a Gaussian filter of 0.4 FWHM. The arrow in the bottom panel shows the direction of the proper motion of PSR B1951+32 (Migliazzo et al. 2002).

Table 1 lists four of them that have isolated counterparts in the Two Micron All Sky Survey (2MASS) point source catalog. The offsets between their Chandra and 2MASS coordinates are smaller than 0.7 arcsec, comparable to the combined systematic uncertainty of the Chandra and 2MASS astrometry. (The systematic uncertainty of the 2MASS astrometry is 0.2 arcsec.)

We observed PSR B1951+32 on 2001 July 16 with the 2048 × 2048 pixel CCD of the Carnegie Observatories Spectrograph and Multi-object Imaging Camera (COSMIC) on the 5 m Palomar Hale telescope. We used g (500 nm), r (655 nm), and i (800 nm) broadband filters with 100 s integration time for each band and proceeded with the standard data reduction, such as bias/dark subtractions and flat-fielding. We next used ~100 2MASS point sources and the IRAF task ccmap to obtain 0.28 arcsec (R.A.) and 0.40 arcsec (decl.) astrometric uncertainties of the COSMIC images. We also performed K band (2.15 μm) observations with the Keck I telescope on 2003 November 6 using the Near-IR Camera (NIRC) equipped with a 256 × 256 pixel InSb detector. We obtained 18 dithered frames, each with 42 s integration time, and subtracted dark and median-combined sky frames. We then shifted and combined them to make a final image. In addition, we reanalyzed archival data for CTB 80 from Wide Field Planetary Camera 2 aboard HST on 1997 October 2. We used ~10 2MASS point sources to obtain the astrometric solutions for the NIRC and HST images with IRAF ccmap.

Figure 2 compares the Hα (HST F656N band; 2.2 nm wide at 656.4 nm) image of the CTB 80 core with the X-ray image. In order to match them, we undertook the following procedure. First, we tied the Chandra image to the COSMIC images using the point sources in Table 1, which gave 0.25 arcsec uncertainty in each coordinate (see the following paragraph for a detailed description). Next, we replaced the COSMIC image with the HST F656N-band image. The astrometric discrepancy between them was estimated to be 0.24 arcsec; therefore, the two images in Figure 2 were matched with 0.35 arcsec uncertainty. In the bottom panel, we identify an ~40″ cometary X-ray nebula, elongated along the proper motion of PSR B1951+32 with the pulsar at the head. The tail shows diffuse and...
extended emission in the northeast, where the pulsar has traveled. The southwestern boundary of the nebula appears to be confined by a bow shock–like feature in Hα, and the X-ray emission shows a steep gradient in the confined region. Figure 3 presents the distribution of the X-ray and Hα emission along the proper motion of PSR B1951+32 (i.e., the arrow in the bottom panel of Fig. 2), showing that the Hα emission peaks $\sim$7 outside of the pulsar in the direction of its proper motion.

Figure 4 shows the COSMIC gri-, HST F547M- (48.6 nm wide at 547.9 nm), and NIRC $K_s$-band images centered on PSR B1951+32. In the COSMIC images, we determined the pulsar position in two different ways. First, we simply found its Chandra position in the COSMIC astrometry with $0.53$ (R.A.) and $0.60$ (decl.) uncertainties, dominated by the systematic astrometric uncertainties of Chandra and 2MASS and by uncertainties in the COSMIC astrometry. The large dotted circle in Figure 4 represents the position of PSR B1951+32 determined in this way. Second, the Chandra image was tied to the COSMIC images using the X-ray point sources and their 2MASS counterparts in Table 1, where the (intrinsic) X-ray positional uncertainties of the point sources and the uncertainties of the COSMIC astrometry dominate the final uncertainties of $0.25$ in each coordinate. The small solid circle in Figure 4 represents the position of PSR B1951+32 determined in this way. The positions of the two claimed optical counterpart candidates (i.e., 1HST and 4HST; Butler et al. 2002) are shown in Figure 4—while 1HST is within the error circle, 4HST is outside.

The optical synchrotron knot of PSR B1951+32 reported by Hester (2000) was identified in the HST F547M-band image in Figure 4e. Given $E(B-V) = 0.8$ extinction toward PSR B1951+32 (Hester & Kulkarni 1989), the extinction-corrected magnitude of the knot is $20.1 \pm 0.2$ [$=(3.3 \pm 0.7) \times 10^{-3}$ Jy] in the F547M band. The knot was also clearly detected in the r and K bands (Figs. 4b and 4d), with estimated extinction-corrected magnitudes of $20.0 \pm 0.2$ [$=(4.5 \pm 1.4) \times 10^{-4}$ Jy] and $18.1 \pm 0.2$ [$=(3.9 \pm 0.7) \times 10^{-3}$ Jy], respectively. (The quoted uncertainties represent the 68.3% confidence levels.) However, emission associated with the knot was not detected in the HST data obtained with narrowband line filters of F502N (2.7 nm wide at 502.3 nm), F656N, and F673N (4.7 nm wide at 673.2 nm).

3. DISCUSSION AND CONCLUSIONS

Previous optical and X-ray observations showed that PSR B1951+32 has likely been interacting with the recombining material behind the radiative shock of the CTB 80 SNR, likely forming a bow shock in the core (e.g., Hester & Kulkarni 1989; Safi-Harb et al. 1995). The radio continuum images resemble our Chandra image, together with a feature reminiscent of a bow shock (e.g., Migliazzo et al. 2002). In Figures 2 and 3, the Hα emission overlaps the X-ray emission in the core (especially in the eastern part), and, around PSR B1951+32, it also shows a bow shock–like feature confining the cometary X-ray nebula.

One simple interpretation of the overall emission of the CTB 80 core (except for the bow shock–associated features around PSR
B1951+32; see below) is that the X-ray and radio emission represent the synchrotron radiation of relativistic pulsar winds, while the Hα emission represents the cooling, recombining thermal plasma shocked by them (e.g., Hester & Kulkarni 1989). Note that this still allows for the existence of Hα emission purely associated with the CTB 80 SNR, which might be responsible for the Hα emission in the western lobes of the CTB 80 core apparently lacking the X-ray and/or radio counterparts, which is quite different from the Hα emission in the eastern part.

Obviously, the most distinctive new feature of the CTB 80 core is the cometary X-ray nebula headed by PSR B1951+32 along its proper motion, seemingly confined by the Hα emission forming bow shock morphology at ∼7″ ahead of the pulsar (Fig. 2). The revelation of such a feature is in good accordance with the previous interpretation of the CTB 80 core (e.g., Hester & Kulkarni 1989), in which the cometary X-ray nebula represents the shocked pulsar winds confined by an Hα bow shock formed by collisional excitation of the ambient medium via the supersonic motion of PSR B1951+32. For this, it is important to note that there is a significant contribution of collisional excitation to the Hα emission in the CTB 80 core (in addition to the recombination mentioned above; Hester & Kulkarni 1989), which is supportive of the bow shock interpretation. One interesting result is that the distance (from the pulsar) to the Hα bow shock (∼7″) is larger than the value (∼3″) obtained for the radio bow shock (Migliazzo et al. 2002; Chatterjee & Cordes 2002), indicating the location of a contact discontinuity for the radio bow shock (Migliazzo et al. 2002; Chatterjee & Kulkarni 2002), which is supportive of the bow shock interpretation.

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In order to form the bow shock, ram pressure balance is required between the relativistic pulsar winds and the ambient medium: $\rho_P v_W^2 = \dot{E}/4\pi \Omega r_s^2$, where $\rho_P$ is the ambient medium density, $v_W$ the pulsar velocity, $\dot{E}$ the pulsar spin-down energy, $4\pi \Omega$ the solid angle through which the pulsar winds flow, $c$ the speed of the light, and $r_s$ the stagnation radius. The assumption of equipartition between the pressure of the magnetic field and the pressure of the flow near the bow shock leads to $B_{\perp}^2 (\mu G) \sim 50 (n_{H_1} \text{cm}^{-3})^{-1} (v_W/100 \text{ km s}^{-1})^2$, where $n_{H_1}$ is the hydrogen number density, and also to $B_{\perp}^2 (\mu G) \sim 20 \Omega^{1/2} (E/10^{36} \text{ ergs s}^{-1})^{1/2} (r/0.01 \text{ pc})^{-1}$. With the observed values $v_W = 240 \text{ km s}^{-1}$, $E = 3.7 \times 10^{36} \text{ ergs s}^{-1}$, and $r_s = 0.03 \text{ pc} (=3')$ of PSR B1951+32, as well as $\Omega < 1$, we estimate that $B_{\perp} > 100 \mu G$. This reconciles with Hester & Kulkarni (1989), who estimated the preshock density and magnetic field to be ∼50 cm$^{-3}$ and ∼600 $\mu$G, respectively. For $B_{\perp} = 600 \mu G$, $\Omega < 0.1$, corresponding to highly anisotropic pulsar winds. If we assume that $B > 100 \mu G$ in the entire cometary X-ray nebula, its synchrotron cooling time is $t_{\text{syn}} (\text{yr}) \sim 40 E_{50}^{1/2} (B/100 \mu G)^{-3/2} < 40 \text{ yr}$, where $E_{50}$ is the observed photon energy in units of keV. Considering the size (∼20″) of the cometary X-ray nebula and the magnitude of the PSR B1951+32 proper motion (∼25 mas yr$^{-1}$), the short cooling time indicates that the cometary X-ray nebula has likely been replenished by the relativistic pulsar winds flowing from the pulsar. Note that similar results have recently been reported for other pulsar wind nebulae (Kaspi et al. 2001; Gaensler et al. 2003). In conclusion, after the millisecond (recycled) pulsar PSR B1957+20 (Stappers et al. 2003), PSR B1951+32 is only the second pulsar showing both inner and outer shocks in a pulsar wind nebula, and it is unique in exhibiting such a feature in the optical, X-ray, and radio emission together.

Our results confirm that the claimed optical synchrotron knot of PSR B1951+32 is indeed a continuum source in nature. In Figure 4, the emission from the knot appears to be present in all five broad bands (although its significance is weak in the $g$ and $i$ bands). For the $HST$ data, the knot is visible only in the relatively line-free F547M-band image (Fig. 4e), while it is absent in the narrowband line filter images (i.e., F502N, F656N, and F673N). Thus, the knot exhibits continuum emission with a flat spectrum in the optical and near-IR wave bands (see § 2 for the flux estimation), unless it has a significant variability. The Crab pulsar is also known to have a similar optical synchrotron knot (Hester et al. 1995) that might be caused by quasi-stationary shocks from pulsar polar outflows (Lou 1998). On the other hand, of the two optical counterpart candidates for PSR B1951+32 (Butler et al. 2002), our improved astrometry is consistent only with 1HST (while it excludes 4HST; Fig. 4). However, it is possible that 1HST simply represents nonuniformity in the optical synchrotron knot or emission from any background star. We need multiepoch, multicolor observations to study the synchrotron knot and optical counterpart more thoroughly.

Part of the data presented herein was obtained at the W. M. Keck Observatory, which is operated by the California Institute of Technology (CIT), the University of California, and NASA. This research is based on the data from the archive at STScI, which is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5-26555. This research has also made use of the data products from the Two Micron All Sky Survey, which is a joint project of the University of Massachusetts and IPAC/CIT, funded by NASA and the NSF. D.-S. M. was supported in part by NSF grant AST-9986898 and also by a Millikan fellowship at CIT; J.-J. L. was supported by KOSEF grant ABRL 3345-20031017; S. S. E. is supported by NSF CAREER award AST-0328522; S. C. is supported by a Jansky fellowship from NRAO; D. L. K. is supported by a fellowship from the Fannie and John Hertz Foundation; Y. A. G. was supported in part by EC fellowship HPMFC-2000-00671.

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