Stark shift and field ionization of arsenic donors in $^{28}$Si-SOI structures

C. C. Lo, S. Simmons, R. Lo Nardo, C. D. Weis, A. M. Tyryshkin, J. Meijer, D. Rogalla, S. A. Lyon, J. Bokor, T. Schenkel, and J. J. L. Morton

1) London Centre for Nanotechnology, University College London, London WC1H 0AH, U.K.
2) Department of Electronic and Electrical Engineering, University College London, London WC1E 7JE, U.K.
3) Department of Materials, University of Oxford, Oxford OX1 3PH, U.K.
4) Accelerator and Fusion Research Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, U.S.A.
5) Department of Electrical Engineering, Princeton University, Princeton, New Jersey 08544, U.S.A.
6) RUBION, Ruhr-Universitaet Bochum, Germany
7) Department of Electrical Engineering and Computer Sciences, University of California, Berkeley, California 94720, U.S.A.

(Dated: 28 January 2014)

We develop an efficient back gate for silicon-on-insulator (SOI) devices operating at cryogenic temperatures, and measure the quadratic hyperfine Stark shift parameter of arsenic donors in isotopically purified $^{28}$Si-SOI layers using such structures. The back gate is implemented using MeV ion implantation through the SOI layer forming a metallic electrode in the handle wafer, enabling large and uniform electric fields up to $\sim 2 V/\mu m$ to be applied across the SOI layer. Utilizing this structure we measure the Stark shift parameters of arsenic donors embedded in the $^{28}$Si SOI layer and find a contact hyperfine Stark parameter of $\eta_c = -1.9 \pm 0.2 \times 10^{-3} \mu m^2/V^2$. We also demonstrate electric-field driven dopant ionization in the SOI device layer, measured by electron spin resonance.

Shallow donors in silicon are promising candidates for spin-based quantum information processing. Although it has been known for decades that the coherence times of donor spins are relatively long at cryogenic temperatures, recent measurements with improved instrumentation, better understanding of the decoherence mechanisms and the availability of ultra high purity isotopically enriched nuclear spin-free $^{28}$Si crystals have led to dramatic improvements of the measured spin coherence times, exceeding seconds for phosphorus and bismuth donor electron spins. Furthermore, the nuclear spin coherence times of ionized phosphorus donors have been found to exceed hours at cryogenic temperatures and 39 minutes at room temperature. In addition, the maximum electric field that could be applied was below $1 V/\mu m$, limited by leakage currents between the electrodes in high-field regions. A parallel-plate capacitor-like structure in principle allows the application of uniform electric fields across the silicon and reduces the probability of having leakage currents. However, the voltage required to bias across the silicon substrate would need to be in the order of 100 $V$ to obtain appreciable Stark shifts for typical wafer thicknesses of 500 $\mu m$.

The first measurement of the Stark shift parameters for shallow donors was reported using electron spin resonance (ESR) of $^{121}$Sb donors residing between interdigitated metallic electrodes evaporated on the silicon surface. The geometry resulted in a non-uniform distribution of the electric field for the donor ensemble. In addition, the maximum electric field that could be applied was $1 V/\mu m$, limited by leakage currents between the electrodes in high-field regions. A parallel-plate capacitor-like structure in principle allows the application of uniform electric fields across the silicon and reduces the probability of having leakage currents. However, the voltage required to bias across the silicon substrate would need to be in the order of 100 $V$ to obtain appreciable Stark shifts for typical wafer thicknesses of 500 $\mu m$.

In this work we develop devices for ensemble measurements utilizing silicon-on-insulator (SOI) substrates, where the donors of interest are embedded in the thin SOI layer. We implement a novel back gate using high energy ion implantation into the handle wafer, which allows the efficient application of uniform and large electric fields across the SOI layer for both Stark shift measurements and controlled ionization of the dopants. This provides

---

a) Electronic mail: cheuk.lo@ucl.ac.uk.

b) Present address: Division of Nuclear Solid State Physics, University of Leipzig, Leipzig, Germany

c) Electronic mail: jjl.morton@ucl.ac.uk.
a platform for ensemble spin resonance measurements on electrically ionized and re-neutralized donor nuclear spins.

The SOI substrates consisted of 100 nm thick p-doped (≈ 10 Ωcm) natural silicon (nat-Si) with \langle 100 \rangle orientation (SOITEC), and a similarly p-doped (≈ 15 Ωcm) handle wafer underneath which normally freezes out at cryogenic temperatures. 100 nm thick isotopically purified 99.92% 28-silicon (28Si) was epitaxially grown on the natural silicon SOI wafer (Lawrence Semiconductors), forming the 28Si-SOI device layer. In order to create an effective back-gate for realizing a parallel plate geometry, a high energy boron (B) implantation of 4 MeV at a dose of 2.2 × 10^{15} /cm^2 was performed, with almost all dopants implanted through the device and the 200 nm thick buried oxide (BOX) layers, peaking at 5 μm down from the handle wafer-buried oxide interface. The peak B concentration of 3 × 10^{19} /cm^3 is above the metal-insulator-transition, ensuring that the p+ region in the handle wafers remains metallic even at cryogenic temperatures. The added B concentration in the SOI layer resulting from the high energy implant is expected to be less than 1 × 10^{15} /cm^3 from device process simulations, and the lattice damage in the SOI layer is expected to be minimal due to ion channeling and most lattice damage is expected to be incurred at end-of-range^{14}. A 10 nm silicon dioxide was subsequently grown on the SOI layer by dry oxidation at 900°C. This oxidation step activated the implanted B and also reduced lattice damage in the handle wafer induced by the MeV implants.

The 28Si-SOI layer was then implanted with arsenic dopants (75As) at 110 keV at a dose of 4 × 10^{11} /cm^2 with 7° tilt. Arsenic was chosen to minimize dopant diffusion from the 28Si layer to the natural silicon layer during thermal activation. In addition, the As nuclear spin of I = 3/2 allows the clear elucidation of the contribution from the contact hyperfine induced Stark shift. The resulting peak donor concentration is located approximately 60 nm away from the dry oxide-SOI layer interface. 360 nm of low temperature oxide was subsequently deposited at 450°C to reduce the probability of gate leakage in the large-area device required for ensemble measurements. The dopants were then activated by rapid thermal annealing at 900°C for 10 s in a nitrogen ambient. We found significant dopant pile up towards the dry oxide as measured by secondary ion mass spectroscopy, which shows an As concentration of 4 – 5 × 10^{10} /cm^3 in the top 50 nm of the 28Si epitaxial layer but drops to ≈ 1 × 10^{15} /cm^3 at the nat-Si-BOX interface, hence 90% of the implanted donors reside in the 28Si layer.

In order to access the deeply implanted p+ back gate in the handle wafer, contact windows were made by etching through the oxide layers and SOI layer. A BOSCH etch process was then carried out to etch 5 μm down the handle wafer. After native oxide removal by hydrofluoric acid, Ti/Au metallic layers were subsequently deposited to form both the contacts to the p+ back gate and defining the top gate electrode. The top gate electrode covers an area of 21 × 2.6 mm^2 with approximately 2 × 10^{11} donors underneath. A forming gas anneal (90% N₂, 10% H₂) at 350°C for 20 minutes was then carried out to passivate the silicon-silicon dioxide interfaces in the structure. A schematic and the cross section of the final device structure are shown in Fig. 1.

Capacitance-voltage (CV) measurements at cryogenic temperatures reveal a gate capacitance of 1.9 nF/cm², in good agreement with the geometry of the device with the top 5 μm of the silicon handle wafer being insulating. The capacitance is essentially constant within the gate voltage (≤ 2 V) and pulse speed (10 μs) used in this study. Frequency dependent CV measurements indicate that hole accumulation at the handle wafer-BOX interface was only achieved at low frequencies (<1 kHz), and on some devices only at temperatures greater than the carrier freeze-out temperature. We speculate that this is due to residual lattice defects in the handle wafer with sufficient trap density to prevent holes from the p+ back gate migrating to the interface efficiently. Hence, the electric field strength is calculated assuming a metallic parallel-plate capacitor geometry, with dimensions as shown in the inset of Fig. 1.

All ESR measurements were carried out with a pulsed X-band (≈ 9.8 GHz and \textit{B} ≈ 0.34 T) ESR spectrometer (Bruker Elexsys 580). A continuous wave ESR linewidth of 0.01 mT (0.3 MHz) for the As donor was measured at 20 K, a factor of ~40 narrower than the known linewidth in nat-Si. In addition, an electron spin coherence time \( T_2 = 0.3 \text{ ms} \) was observed at 2 K, indicating that the spin coherence is limited by the relatively high As concentration instead of nat-Si in the SOI layer.\(^\text{5}\) Both of these observations confirm that the signal mostly arises from dopants residing in the 28Si layer and that the SOI layer has a low defect density. All measurements discussed hereafter were performed at 15 K, unless stated otherwise.

FIG. 1. Schematic of the 28Si-SOI device with a p+ back gate formed by MeV ion implantation. Inset: Cross-section of the device. 90% of the 75As donors reside in the epitaxially grown 28Si region of the SOI layer.
Our parallel plate structure permits a uniform electric field to be readily applied to the As donors. Under small applied voltages the donor remains neutral at low temperatures, however, the electronic wavefunctions are distorted, causing a change in the contact hyperfine interaction and spin-orbit coupling strengths, hence inducing a Stark shift. We investigate the As Stark shift parameters using a Hahn echo pulse sequence, and we detect the frequency shift of the donors by examining the change in the Hahn echo signal phase as the applied voltage (electric field) is varied. The static magnetic field $B_0$ is aligned along the (110) axis, and orthogonal to the direction of the applied DC electric field $E$.

We use a bipolar voltage pulse sequence (sequence IV in Ref [13]), where voltages of opposite sign but equal amplitude are applied successively during the defocusing phase of the Hahn echo sequence. The synchronization of the microwave and voltage pulses is shown in Fig. 2(a). The advantage of using a bipolar pulse sequence over a unipolar pulse sequence (containing only voltage pulses of the same sign) is that ambiguities arising from built-in electric fields cancel out. The donor resonant frequency shift $\Delta f_{m_j}$ for a given ESR line corresponding to nuclear spin $m_j$, due to the applied electric field of magnitude $|E|$ is given by:

$$\Delta f_{m_j} = \left( \frac{1}{\hbar} \right) (\eta_a a m_j + \eta_g g \mu_B B) E^2 \quad (1)$$

where $\eta_a$ is the hyperfine induced Stark parameter, $\eta_g$ the spin-orbit induced Stark parameter, $\hbar$ Planck’s constant and $\mu_B$ the Bohr magneton, while $a$ and $g$ are the hyperfine interaction strength and electronic $g$-factor, respectively. For As donors, $a = 819$ neV and $g = 1.99837^1$. We have neglected linear Stark shift terms ($\propto E$) which are much smaller compared with the quadratic effects in silicon$^{14,15}$, and linear terms cancel out with the bipolar voltages pulses used. With such a frequency shift, a voltage (electric field) pulse of duration $\tau$ then induces a phase shift $\Delta \varphi_{m_j} = 2\pi \Delta f_{m_j} \tau$ for the Hahn echo signal.

All measurements were carried out with microwave pulse delays of $\tau_0 = 15 \mu s$, and the voltage pulses for each polarity are $5 \mu s$ long (total $10 \mu s$ voltage pulse for each measurement). A two step phase cycling (+X,+X and +X,−X) for the microwave pulses is used to remove spurious background signals. Figure 2(b) shows the reference echo signal and Fig. 2(c) shows the phase shifted echo signal under an applied electric field of $E = 0.23 \, V/\mu m$ for the $m_j=+3/2$ line. The echo phase shifts extracted from the echo signals for the four different hyperfine-split lines are converted to frequency shifts and shown in Fig. 2(d). The solid lines represent quadratic fits to the data as defined by Eq. 1. We find the contact hyperfine term $\eta_a = -1.9 \pm 0.2 \times 10^{-3} \, \mu m^2/V^2$ and the spin-orbit term $\eta_g = 3 \pm 4 \times 10^{-6} \, \mu m^2/V^2$. The spin-orbit shift is very weak compared with the hyperfine shift under our experimental conditions with relatively moderate magnetic fields $^{13,16}$, and cannot be determined accurately within experimental errors. Indeed, fitting the data while neglecting the spin-orbit term results in virtually no change to the contact hyperfine value. We summarize the Stark shift parameters for various silicon shallow donors in Table I.

We find that at larger electric fields the ESR signal

![Figure 2](image_url)
we found that even at 1 V/µm the ESR signals are significantly reduced, indicating that some donors begin to be ionized by the moderate electric fields, which is consistent with theoretical calculations and electrical measurements. This underlines the difficulty of Stark tuning of donors relying on the hyperfine term alone. At higher magnetic fields when the spin-orbit term becomes appreciably larger, or by inducing additional strain in the sample, larger Stark shifts can be induced to circumvent this issue.

In summary, we have investigated donor spin dynamics under the influence of applied electric fields with pulsed electron spin resonance. Arsenic donors are embedded in isotopically purified 28-silicon SOI parallel-plate structures formed between a metallic top gate and MeV implanted back gate. We measured the Stark shift of the donors by applying electric fields below the ionization threshold of the dopants, and find a hyperfine Stark shift parameter of $\eta_a = -1.9 \pm 0.2 \times 10^{-3}$ µm²/V².

### ACKNOWLEDGMENTS

This research is supported by the EPSRC through the Materials World Network (EP/I035536/1) and CAESR (EP/D048559/1) as well as by the European Research Council under the European Community’s Seventh Framework Programme (FP7/2007-2013)/ERC grant agreement No. 279781. Work at Princeton was supported by the National Security Agency (100000080295) and by the U.S. Department of Energy (DE-AC02-05CH11231, LBNL). C.C.L. is supported by the Royal Commission for the Exhibition of 1851. S.S. is supported by the Violette and Samuel Glasstone Fund and St. Johns College Oxford. J.J.L.M. is supported by the Royal Society.

---

**TABLE I. Summary of the experimentally and theoretically determined Stark shift parameters for shallow donor in silicon.**

| Donor Orientation | $\eta_a$ | $\eta_b$ | Experiment/Theory |
|-------------------|----------|----------|-------------------|
| $^{31}$P          | –        | -2.8     | Tight-binding calculation.¹⁵ |
| $B \parallel E$   | –        | -12      | Tight-binding calculation.¹⁶ |
| $B \perp E$       | –        | +14      | Tight-binding calculation.¹⁶ |
| $^{75}$As         | $B \perp E$ | $-1.9 \pm 0.2$ | $+3 \pm 4$ | ESR with parallel-plate geometry (present work). |
| $^{121}$Sb       | $B \perp E$ | $-3.7$   | $-10$ | ESR with interdigitated electrodes.¹³⁸ |

Note that the signs of the nuclear spin projections in Fig. 4 of Ref.¹⁵ were mislabeled, but the signs and values given for the Stark shift parameters are correct.

---

As dopant density used in this study corresponds to a mean donor-donor separation of approximately 40 nm, which is close to the upper limit of exchange-based entanglement schemes for donor qubits. At these concentration levels, the spin coherence time is limited by direct dipole-dipole interaction amongst the donors.⁵

We note that in a perfectly planer device architecture the dipole-dipole interaction can be removed by applying the magnetic field at the magic angle.⁷ However, the line widths are still dominated by inhomogeneous broadening, which is about 0.2 MHz in the highest quality $^{28}$Si material, implying an electric field of the order of 1 V/µm is needed for shifting an entire ESR line out of resonance as previously noted.¹⁵ We found that even at 1 V/µm the ESR signals are significantly reduced, indicating that some donors begin to be ionized by the moderate electric fields, which is consistent with theoretical calculations and electrical measurements. This underlines the difficulty of Stark tuning of donors relying on the hyperfine term alone. At higher magnetic fields when the spin-orbit term becomes appreciably larger, or by inducing additional strain in the sample, larger Stark shifts can be induced to circumvent this issue.

In summary, we have investigated donor spin dynamics under the influence of applied electric fields with pulsed electron spin resonance. Arsenic donors are embedded in isotopically purified 28-silicon SOI parallel-plate structures formed between a metallic top gate and MeV implanted back gate. We measured the Stark shift of the donors by applying electric fields below the ionization threshold of the dopants, and find a hyperfine Stark shift parameter of $\eta_a = -1.9 \pm 0.2 \times 10^{-3}$ µm²/V².
2J. P. Gordon and K. D. Bowers. Microwave spin echoes from donor electrons in silicon. Physical Review Letters, 1(10):368, 1958.

3G. Feher. Electron spin resonance experiments on donors in silicon. I. Electronic structure of donors by electron nuclear double resonance technique. Physical Review, 114(5):1219, 1959.

4T. G. Castner. Orbach spin-lattice relaxation of shallow donors in silicon. Physical Review, 155(3):154, 1967.

5A. M. Tyrnyshkin, S. Tojo, J. J. L. Morton, H. Riemann, N. V. Abrosimov, P. Becker, H. J. Pohl, T. Schenkel, M. L. W. Thewalt, K. M. Itoh, and S. A. Lyon. Electron spin coherence exceeding seconds in high-purity silicon. Nature Materials, 11(2):143, 2012.

6G. Wolfowicz, A. M. Tyrnyshkin, R. E. George, H. Riemann, N. V. Abrosimov, Peter Becker, Hans-Joachim Pohl, M. L. W. Thewalt, S. A. Lyon, and John J. L. Morton. Atomic clock transitions in silicon-based spin qubits. Nature Nanotechnology, 8:561, 2013.

7K. Saeedi, S. Simmons, J. J. L. Morton, and M. L. Thewalt. Room-temperature quantum bit storage exceeding 39 minutes using ionized donors in silicon-28. Science, 342:830, 2013.

8A. Morello, J. J. Pla, F. A. Zwanenburg, K. W. Chan, K. Y. Tan, H. Huebl, M. Möttönen, C. D. Nugroho, C. Yang, J. A. van Donkelaar, A. D. C. Alves, D. N. Jamieson, C. C. Escott, L. C. L. Hollenberg, R. G. Clark, and A. S. Dzurak. Single-shot readout of an electron spin in silicon. Nature, 467:687, 2010.

9J. J. Pla, K. Y. Tan, J. P. Dehollain, W. H. Lim, J. J. L. Morton, D. N. Jamieson, A. S. Dzurak, and A. Morello. A single-atom electron spin qubit in silicon. Nature, 489:541, 2012.

10J. J. Pla, K. Y. Tan, J. P. Dehollain, W. H. Lim, J. J. L. Morton, F. A. Zwanenburg, D. N. Jamieson, A. S. Dzurak, and A. Morello. High-fidelity readout and control of a nuclear spin qubit in silicon. Nature, 496:334, 2013.

11F. A. Mohiyaddin, R. Rahman, R. Kaira, G. Klimeck, L. C. L. Hollenberg, J. J. Pla, A. S. Dzurak, and A. Morello. Noninvasive spatial metrology of single-atom devices. Nano Letters, 13:1903, 2013.

12L. Dreher, F. Hoehne, M. Stutzmann, and M. S. Brandt. Nuclear spins of ionized phosphorus donors in silicon. Physical Review Letters, 108(2):027602, 2012.

13F. R. Bradbury, A. M. Tyrnyshkin, G. Sabouret, J. Bokor, T. Schenkel, and S. A. Lyon. Stark tuning of donor electron spins in silicon. Physical Review Letters, 97(17):176404, 2006.

14SRIM simulator. The Stopping and Range of Ions in Matter. http://www.srim.org/, 2010.

15R. Rahman, C. J. Wellard, F. R. Bradbury, M. Prada, J. H. Cole, G. Klimeck, and L. C. L. Hollenberg. High precision quantum control of single donor spins in silicon. Physical Review Letters, 99(3):036403, 2007.

16R. Rahman, S. H. Park, T. B. Boykin, G. Klimeck, S. Rogge, and L. C. L. Hollenberg. Gate-induced g-factor control and dimensional transition for donors in multivalley semiconductors. Physical Review B, 80(15):155301, 2009.

17R. de Sousa and J. D. Delgado S. D. Sarma. Silicon quantum computation based on magnetic dipolar coupling. Physical Review A, 70(5):052304, 2004.

18M. J. Calderon, B. Koiller, and S. Das Sarma. External field control of donor electron exchange at the Si/SiO₂ interface. Physical Review B, 75(12):125311, 2007.

19S. Zurauskas and A. Dargys. Field ionization of phosphorus atoms in silicon. Phys. Stat. Sol. (b), 121:385, 1984.

20L. Dreher, T. A. Hilker, A. Brandlmaier, S. T. B. Goennenwein, H. Huebl, M. Stutzmann, and M. S. Brandt. Electrostatic hyperfine tuning of phosphorus donors in silicon. Physical Review Letters, 106(3):037601, 2011.