THE IMPLICATIONS OF EXTREME OUTFLOWS FROM EXTREME STARBURSTS

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ABSTRACT
Interstellar ultraviolet absorption lines provide crucial information about the properties of galactic outflows. In this paper, we augment our previous analysis of the systematic properties of starburst-driven galactic outflows by expanding our sample to include a rare population of starbursts with exceptionally high outflow velocities. In principle, these could be a qualitatively different phenomenon from more typical outflows. However, we find that instead these starbursts lie on, or along the extrapolation of, the trends defined by the more typical systems studied previously by us. We exploit the wide dynamic range provided by this new sample to determine scaling relations of outflow velocity with galaxy stellar mass ($M_*$), circular velocity, star formation rate (SFR), SFR/$M_*$, and SFR/area. We argue that these results can be accommodated within the general interpretational framework we previously advocated, in which a population of ambient interstellar or circumgalactic clouds is accelerated by the combined forces of gravity and the momentum flux from the starburst. We show that this simple physical picture is consistent with both the strong cosmological evolution of galactic outflows in typical star-forming galaxies and the paucity of such galaxies with spectra showing inflows. We also present simple parameterizations of these results that can be implemented in theoretical models and numerical simulations of galaxy evolution.

Key words: galaxies: evolution – galaxies: ISM – galaxies: kinematics and dynamics – galaxies: starburst – intergalactic medium

1. INTRODUCTION
Galactic outflows driven by the energy and momentum supplied by a population of short-lived massive stars play crucial roles in the evolution of galaxies and the intergalactic medium (Somerville & Davé 2015 and references therein). They can provide feedback that expels existing gas-phase baryons from galaxies and can prevent or inhibit the accretion of new gas from the circumgalactic or intergalactic medium. This feedback can help explain the small ratio of baryons to dark matter observed in low-mass galaxies (e.g., McGaugh et al. 2010). In addition, the selective loss of newly synthesized heavy elements from shallow potential wells is a key factor in the galaxy mass–metallicity relation (e.g., Tremonti et al. 2004; Andrews & Martini 2013). The outward transport of low-angular momentum gas in outflows may help govern the relationship between the mass and size of galactic disks (e.g., Kauffmann et al. 2003). Galactic outflows have also been effective at transporting metals and even dust out into the circumgalactic and intergalactic mediums (e.g., Ménard et al. 2010).

While there is general qualitative agreement about the importance of galactic outflows, it has not been possible so far to reliably quantify their effects. Part of the problem is that numerical simulations of galaxy evolution in a cosmological context cannot incorporate the physics of galactic outflows in a fully ab initio manner. Instead, both simulations and semi-analytic models are generally forced to rely on simple parameterizations of the relevant processes (Somerville & Davé 2015). Ideally, these would be based on a robust empirical characterization and theoretical understanding of galactic outflows, but neither of these yet exist.

One of the chief tools for investigating galactic outflows is spectroscopy using resonance lines in absorption to probe the cool and warm phases of the outflow (e.g., Heckman et al. 2000; Steidel et al. 2010). This has the advantage that it can be readily applied to star-forming galaxies at both low-redshift (where our observational characterization of outflows is most complete) and high-redshifts (where galactic outflows are ubiquitous and strong).

Motivated by these considerations, we ( Heckman et al. 2015, hereafter H15) recently analyzed high-quality ultraviolet spectroscopic data for a sample of 39 low-redshift starburst galaxies that spanned broad ranges in the principal properties of the galaxies and their starbursts. We found that some (and some combination) of these properties had strong systematic correlations with the observed properties of the outflows. We also showed that a simple model in which the outflowing gas seen in absorption is produced by a population of clouds accelerated by a combination of gravity and the momentum flux supplied by the starburst provided a good fit to the data.

In this paper, we seek to test these correlations and our model by extending our analysis to a population of extreme starbursts that significantly extends the range in outflow velocity compared to the sample in H15. In principle, these high outflow velocities might require a qualitatively different model than that advanced in H15. For example these could be outflows driven by a powerful active galactic nucleus (AGN) (e.g., Liu et al. 2013) that has recently shut-down (e.g., LaMassa et al. 2015) or thermal instabilities in a radiatively cooling fast outflow of thermalized supernova and stellar wind ejecta (Thompson et al. 2016).

We describe the new expanded sample in Section 2 below, we highlight the empirical correlations in Section 3, and briefly comment on the implications of these results in Section 4.

2. SAMPLE PROPERTIES
In H15 we investigated a sample of 39 low-redshift ($z < 0.2$) starburst galaxies using observations of their ultraviolet interstellar absorption lines made with the Far Ultraviolet Spectroscopic Explore (Grimes et al. 2009) and with the Cosmic Origins Spectrograph on the Hubble Space Telescope (Alexandroff et al. 2015). These papers contain full
descriptions of both the sample selection and the data analysis, and we refer the reader to them for details. In brief, about 60% of the sample consisted of Lyman Break Analogs, which have properties very similar to those of typical Lyman Break Galaxies at redshifts z ~ 3–4 (e.g., Hoopes et al. 2007; Overzier et al. 2010). The other 40% of the sample consisted of ultraviolet-light-low-redshift galaxies representative of the local starburst population.

In this paper, we add a new sample of extreme starbursts drawn from the recent investigations by Diamond-Stanic et al. (2012) and Sell et al. (2014; hereafter S14). These are intermediate redshift (z ~ 0.4–0.7) galaxies discovered in the Sloan Digital Sky Survey archive, characterized by very high outflow velocities (over ~10^3 km s^{-1}), high star formation rates (SFRs) (a few hundred M_☉ yr^{-1}) and compact sizes (starburst half-light radii of a few hundred pc). These values lie well beyond the ranges covered by the H15 sample.

Based on their multi-waveband analysis, S14 concluded that the majority of these objects are energetically dominated by the intense and compact starburst, with little or no contribution by an AGN. From their full sample of 12 galaxies, we have excluded the four that show any evidence for an AGN. We have also added the very similar galaxy investigated by Geach et al. (2014; hereafter G14), giving us a sample of nine extreme starbursts.

While there have been a number of other investigations of outflows driven from intermediate redshift star-forming galaxies (e.g., Weiner et al. 2009; Erb et al. 2012; Kornei et al. 2012; Martin et al. 2012; Bordoloi et al. 2014; Rubin et al. 2014), these galaxies have been representative of typical star-forming galaxies at these epochs. Thus, the extreme-starburst sample complements these earlier studies by extending the investigation of outflows into an entirely new part of parameter space.

The principal properties of the new sample of extreme starbursts are listed in Table 1. The values for galaxy stellar mass (M_g), SFR, starburst half-light radius (r_s), and maximum outflow velocity (v_{max}) are taken directly from S14 and G14. For the most part, these parameters were derived in a way that is consistent with the approach taken for the sample analyzed in H15.

One exception is the definition of the outflow velocity. These were based on the Mg II 2796,2803 doublet for the extreme starbursts, and in many cases the doublet shows a significant amount of line-emission near, and redward of, the galaxy systemic velocity. In this case, in-filling of the absorption-line by emission can have a serious effect on the net profile (e.g., Prochaska et al. 2011; Erb et al. 2012; Kornei et al. 2012; Martin et al. 2012; Rubin et al. 2014; Scarlata & Panagia 2015). This will mean that the flux-weighted centroid of the absorption lines will tend to overestimate the outflow velocity. The maximum blueward extent of the absorption lines will not be affected by infilling, so we adopt these values from S14. We have reanalyzed the COS and FUSE spectra described in H15 and measured the average value of the maximum outflow velocity based on the Si ii 1260 and C ii 1334 lines (COS) and the C II 1036 line (FUSE). These transitions arise from species that roughly match the ionization state of the Mg ii ion. We will refer to these outflow velocities as v_{max} to distinguish them from the flux-weighted mean outflow velocities (v_{out}) investigated in H15.

We have estimated the uncertainty in v_{max} by comparing the values measured individually for Si ii 1190, Si ii 1206, Si ii 1260, and C ii 1334 (COS data) and for C ii 977, C ii 1036, and N ii 1084 (FUSE data). These values are listed in Table 2.

As in H15, we define the SFR/area to be 0.5 SFR/πr_{s}^2 and estimate the galaxy circular velocity (v_{circ}) based on the tight empirical relationship shown by the data presented by Simons et al. (2015) and then parameterized by H15: log v_{circ} = 0.29 log M_g - 0.79, where v_{circ} is in km s^{-1}, and M_g is in solar masses.

3. RESULTS

The motivation of this paper is to extend the analysis of outflows into the regime of extreme starbursts. In Figure 1 we plot a set of correlations between the outflow velocity and the principal properties of the galaxies and their starbursts. It is immediately clear that the extreme starburst sample allows us to probe hitherto unexplored parts of parameter space.

Different investigations in the past have not always agreed with one another in terms of the strengths of the correlations of outflow velocity with SFR (Heckman et al. 2000; Martin 2005; Rupke et al. 2005; Weiner et al. 2009; Erb et al. 2012; Kornei et al. 2012; Martin et al. 2012; Bordoloi et al. 2014; Rubin et al. 2014; Chisholm et al. 2015; H15), with SFR/area (Chen et al. 2010; Kornei et al. 2012; Rubin et al. 2014; Chisholm et al. 2015; H15), and with either M_g or v_{circ} (Heckman et al. 2000; Martin 2005; Rupke et al. 2005; Erb et al. 2012; Martin...
et al. 2012; Bordoloi et al. 2014; Rubin et al. 2014; Chisholm et al. 2015; H15).

The new sample allows us to re-examine the relations between outflow velocity and these properties over an unprecedented dynamic range. These relations are plotted in Figure 1. The first important conclusion is that in all four panels the extreme starbursts lie along an extrapolation of the relations defined by the more typical starbursts in the H15 sample. This provides indirect evidence that these extreme outflows are not a qualitatively different physical phenomenon.

Figure 1 shows that the maximum outflow velocities correlate most strongly with SFR/area, and least strongly with SFR/$M_*$, and $v_{\text{max}}$, we can quantify this by determining simple analytic fitting relations to the correlations in the four panels. The relations with $v_{\text{circ}}$, SFR, and SFR/$M_*$ can all be fit as single power-laws (which we give). The correlation with SFR/area saturates at the high end, and so we fit this as a double power law. We also list the rms residuals about the fits in each panel.

The relationship between $v_{\text{max}}$ and $v_{\text{circ}}$ is particularly important in the context of models of galaxy evolution (Somerville & Davé 2015 and references therein). The ratio of these velocities in Figure 1 varies by about an order-of-
magnitude. We now show in Figure 2 that this ratio correlates strongly and systematically with the star formation rate (SFR/area), and less so with the specific SFR (SFR/$M_*$). The results are consistent with those presented in H15, but the relationships can now be probed over an increased dynamic range (particularly for SFR/area). The correlation with SFR/area also shows a saturation in normalized outflow velocity at $v_{\text{max}} \sim 6-10$ $v_{\text{circ}}$ above SFR/area $\sim 10^2 M_\odot$ yr$^{-1}$ kpc$^{-2}$. It is important to emphasize that outflow velocities this high are well in excess of the galaxy escape velocities. As in the case of Figure 1, we show the analytic fits to the relations and list the rms residuals about these fits.

### 4. DISCUSSION

#### 4.1. The Model of Momentum-driven Clouds

We have shown that the extreme starbursts lie along the extrapolation of the trends defined by the more typical starbursts (Figures 1 and 2). We now want to discuss whether these results can be easily understood within the context of the basic physical picture described in H15.

H15 showed that a simple model of a population of clouds that are accelerated by the combination of gravity and the starburst momentum flux provided a good description of the properties of galactic outflows. In particular, we defined a critical momentum flux such that the net force on a cloud located at $r = r_*$ (the launch point of the outflow) is outward:

$$p_\text{crit} = \Omega r_\ast N_c \langle m \rangle v_{\text{circ}}^2.$$  \hfill (1)

Here, $\Omega$ is the solid angle occupied by the wind, $N_c$ is the cloud Hydrogen column density and $\langle m \rangle$ is the mean mass per H atom. In convenient units $p_\text{crit} = 10^{33.9}$ dynes for $\Omega = 4\pi$, $N_c = 10^{21}$ cm$^{-2}$, $r_\ast = 1$ kpc, and $v_{\text{circ}} = 100$ km s$^{-1}$.

The momentum flux from the starburst is a combination of contributions from both radiation pressure (e.g., Murray et al. 2005) and the ram pressure of a hot outflowing wind fluid collectively created from the ejecta of massive stars (Chevalier & Clegg 1985). For a standard Kroupa/Chabrier initial mass function and a constant SFR

$$p_\ast = 4.8 \times 10^{33} \text{SFR dynes},$$  \hfill (2)

where SFR is in $M_\odot$ yr$^{-1}$ (H15). By expressing the SFR in gm s$^{-1}$ (hereafter, sfr), we can rewrite this as

$$p_\ast = v_{\text{eff}}sfr,$$  \hfill (3)

where $v_{\text{eff}} = 7.6 \times 10^7$ cm s$^{-1}$ (760 km s$^{-1}$).

We showed in H15 that the properties of the outflows depended strongly on the ratio of the momentum flux supplied by the starburst ($p_\ast$) relative to the critical value ($p_\text{crit}$):

$$R_\text{crit} = p_\ast/p_\text{crit} = sfr v_{\text{eff}}/\Omega r_\ast N_c \langle m \rangle v_{\text{circ}}^2.$$  \hfill (4)

In convenient units, $R_\text{crit} = 0.57$ for SFR = 1 $M_\odot$ yr$^{-1}$, $\Omega = 4\pi$, $r_\ast = 1$ kpc, $N_c = 10^{21}$ cm$^{-2}$, and $v_{\text{circ}} = 100$ km s$^{-1}$. We cannot directly compute $R_\text{crit}$ for the extreme starbursts, since we have no estimate of $N_c$ for them. Nonetheless, the very high SFR and small sizes of these objects imply values for $R_\text{crit}$ at least as high of the upper end of the H15 sample ($R_\text{crit} > 10$). Indeed, in Figures 1 and 2 the extreme outflows (red points) overlap much better with the H15 strong outflows (blue points, defined as $R_\text{crit} > 10$) than with the weak outflows (green points, $1 < R_\text{crit} < 10$).

### Table 2

| Galaxy          | $v_{\text{max}}$ (km s$^{-1}$) |
|-----------------|---------------------------------|
| J0021+00        | 350                             |
| J0055-00        | 530                             |
| J0150-13        | 450                             |
| J0213-12        | 1500                            |
| J0808+39        | 1500                            |
| J0823+28        | 370                             |
| J0921+45        | 1500                            |
| J0926+44        | 550                             |
| J0938+54        | 520                             |
| J1025+36        | 360                             |
| J1112+55        | 990                             |
| J1113+29        | 510                             |
| J1144-40        | 570                             |
| J1414+05        | 370                             |
| J1416+12        | 780                             |
| J1428+16        | 440                             |
| J1429+06        | 660                             |
| J1521+07        | 490                             |
| J1525+07        | 700                             |
| J1612+08        | 1000                            |
| J2103-07        | 1260                            |
| Haro 11         | 290                             |
| VV 114          | 400                             |
| NGC 1140        | 150                             |
| SBS 0335-052    | 60                              |
| Tol0400-381     | 230                             |
| NGC 1705        | 170                             |
| NGC 1741        | 190                             |
| I Zw 18         | 90                              |
| NGC 3310        | 630                             |
| Haro 3          | 210                             |
| NGC 3690        | 340                             |
| NGC 4214        | 150                             |
| IRAS 19245+4140 | 210                             |
| NGC 7673        | 230                             |
| NGC 7714        | 380                             |

**Notes:** The ratio of these velocities in Figure 1 varies by about an order-of-magnitude. We now show in Figure 2 that this ratio correlates strongly and systematically with the star formation rate (SFR/area), and less so with the specific SFR (SFR/$M_*$). The results are consistent with those presented in H15, but the relationships can now be probed over an increased dynamic range (particularly for SFR/area). The correlation with SFR/area also shows a saturation in normalized outflow velocity at $v_{\text{max}} \sim 6-10$ $v_{\text{circ}}$ above SFR/area $\sim 10^2 M_\odot$ yr$^{-1}$ kpc$^{-2}$. It is important to emphasize that outflow velocities this high are well in excess of the galaxy escape velocities. As in the case of Figure 1, we show the analytic fits to the relations and list the rms residuals about these fits.
H15 derived the equation-of-motion for the clouds in the model we are considering here. In particular, they showed that the maximum outflow velocity for a cloud in an isothermal potential (which will occur at the radius at which the net radial force on the cloud is zero) could be written as:

\[ v_{\text{max}, \text{cir}} = \frac{v_{\text{circ}}}{R_{\text{crit}}} \sqrt{\frac{2}{R_{\text{crit}}}} \left( R_{\text{crit}} - 1 \right) - \ln(R_{\text{crit}}) \]  

H15 showed that the predicted outflow velocities were a good match to the data. Without direct measurements of \( R_{\text{crit}} \) for the extreme starbursts, we cannot extend this comparison to the extreme starbursts. However, as discussed in H15 Equation (5) above predicts a rapid rise in outflow velocity for small values (\( R_{\text{crit}} \sim 1 \) to 10) and then a flattening in the relationship for \( R_{\text{crit}} > 10 \). Given that \( R_{\text{crit}} \propto \text{SFR}/M_{*} \), this is almost certainly the underlying physics seen in the flattening of the upper end of the relationship between \( v_{\text{max}}/v_{\text{cir}} \) and SFR/area (Figure 2).

We conclude that the very large outflow velocities seen in the extreme starbursts are at least qualitatively consistent with model of momentum-driven clouds in H15 (taken into an extreme regime).

**4.2. Implications**

As we emphasized in Section 1, having a secure empirical characterization and physical understanding of galactic outflows is an important step in being able to develop a quantitative assessment of their impact. We therefore close the paper by considering the implications of our results for understanding galaxy evolution.
The Astrophysical Journal, 822:9 (6pp), 2016 May 1

First, it is instructive to recast \( R_{\text{crit}} \) as defined in H15 and summarized above. We define the dynamical mass of the starburst as

\[
M_{\text{sh}} = 2v_c^2 r_s / G,
\]

where this assumes that half the mass is enclosed within the half-light radius \( r_s \). This together with Equation (5) above then allows us to write:

\[
R_{\text{crit}} = 2s\text{ffr}_{\text{eff}} / \Omega G M_{\text{sh}} N_c \langle m \rangle.
\]

This equation shows that \( R_{\text{crit}} \propto \text{SFR}/M_{\text{sh}} = \text{sFFR}_{\text{sh}} \). We emphasize that this is a specific SFR within the starburst and is normalized with respect to a dynamical mass. In both respects, this is different from the specific SFR pertaining to the entire galaxy and normalized to the total galaxy stellar mass (as plotted in Figures 1 and 2). In convenient units, \( R_{\text{crit}} = 2.7(4\pi/\Omega)(\text{SFR}_{\text{sh}} / \text{Gyr}^{-1}) (N_c/10^{21} \text{ cm}^{-2}) \).

Low-redshift starbursts are usually compact and circumnuclear (i.e., \( M_{\text{sh}} \ll M_* \)). This is not typically the case at high-redshift (e.g., Forster-Schreiber et al. 2011; Shapley 2011; Somerville & Davé 2015). If we therefore make the assumption that the galaxy-wide specific SFR can be used to estimate \( R_{\text{crit}} \) at high-redshift, the strong evolution in the specific SFR with redshift (Madau & Dickinson 2014 and references therein) implies a correspondingly large increase in \( R_{\text{crit}} \) and hence in the importance of strong outflows. More quantitatively, adopting \( N_c = 10^{21} \text{ cm}^{-2} \) would imply that the characteristic value of \( R_{\text{crit}} \) for galaxies on the star-forming main sequence increases from \( \sim 0.3 \) (\( z \sim 0 \)), to \( \sim 3 \) (\( z \sim 1 \)), to \( \sim 5 \) (\( z \sim 2 \)), and to \( \sim 13 \) (\( z \sim 4 \) to 7). This is at least qualitatively consistent with the rarity of strong outflows in the present-day universe and their near- ubiquity at \( z > 2 \).

These results may also have implications for understanding why it has been difficult to find direct spectroscopic evidence for inflowing gas in strongly star-forming galaxies at intermediate and high redshift (e.g., Steidel et al. 2010; Martin et al. 2012; Rubin et al. 2012). To begin, we note that an infalling cloud must satisfy the condition that the inward force of gravity exceeds the outward momentum flux on the cloud from the starburst. This is just recasting Equation (7) above in terms of a critical column density for infall:

\[
N_{\text{inflow}} > N_{\text{crit}} = 2s\text{ffr}_{\text{sh}} v_{\text{eff}} / (\Omega G m). \tag{8}
\]

Here \( s\text{ffr}_{\text{sh}} \) is in units of \( s^{-1} \) and \( v_{\text{eff}} \) is \( 7.6 \times 10^7 \text{ cm s}^{-1} \). In convenient units this corresponds to \( N_{\text{crit}} = 8 \times 10^{21} \text{ cm}^{-2} \) for \( s\text{FFR}_{\text{sh}} = 10^{-9} \text{ yr}^{-1} \) and \( \Omega = 4\pi \). Assuming a normal dust/metal ratio in this material (Mattsson et al. 2014), the Calzetti et al. (2000) dust attenuation law would imply a far-UV extinction of \( A_{\text{fuv}} = 25Z/Z_\odot \) magnitudes for this column density. If the outflow flux is not spherically symmetric (e.g., it is bipolar or more generally follows the path-of-least resistance out of the galaxy) the implied column density and dust extinction become even larger.

Thus, unless the gas has very sub-solar abundances (which might apply to relatively pristine gas being accreted from the cosmic web), it would be essentially opaque in the far-UV and therefore undetectable in the spectra. If this gas covered the whole far-UV source, the galaxy itself would be invisible in the far-UV and would not even enter a sample targeted for rest-frame far-UV spectroscopy in the first place.

5. CONCLUSIONS

The goal of this paper was to improve our understanding of galactic outflows by maximizing the dynamic range over which their properties can be probed. To that end, we have used observations of outflows traced by ultraviolet interstellar absorption lines for a sample of nine extreme starbursts (Geach et al. 2014; Sell et al. 2014). More specifically, the extreme starbursts are characterized by significantly higher star formation rates per unit area (SFR/area) than even the most extreme members of the sample of 39 low-redshift starbursts we studied previously (Heckman et al. 2015, hereafter H15).

We found that in all respects, the extreme starbursts lay along a smooth extrapolation of the correlations seen in H15. The addition of the extreme starbursts strengthened the results...
in H15 by significantly expanding the dynamic range over which these correlations have been delineated.

Specifically, we found that the maximum outflow velocity ($v_{\text{max}}$) correlated most strongly with the SFR/area and least strongly with SFR/$M_\star$. The ratio of $v_{\text{max}}/v_{\text{crit}}$ spanned about an order-of-magnitude and correlated strongly and positively with SFR/area (and less so with SFR/$M_\star$). This ratio reached typical values of $\sim 3 - 10$ for starbursts with high SFR/area, well in excess of the galaxy escape velocity. We exploited the large dynamic range spanned by our sample to derive simple analytic fits to all these empirical relations.

We then argued that the properties of the extreme starbursts were consistent with the simple analytic model for the outflows explored by H15 in which a population of clouds is accelerated by the net sum of gravity and the momentum-flux supplied by the starburst. H15 emphasized the importance of the ratio of the momentum flux supplied by the starburst to the minimum amount required to balance the inward force of gravity on a cloud ($R_{\text{crit}}$).

We showed that $R_{\text{crit}}$ is simply proportional to the value of SFR per unit dynamical mass evaluated over the star-forming region within the galaxy. We argued that the strong observed evolution in the specific SFR with cosmic time then implies a strong evolution in the importance of strong outflows. This is at least qualitatively consistent with what we know observationally. We also showed that material meeting the criterion $R_{\text{crit}} < 1$ (which is required for infall) will have column densities of order $10^{22}$ cm$^{-2}$ at intermediate and high redshift. This material would be opaque in the rest-frame far-UV, potentially explaining why the direct signature of infalling gas (redshifted absorption lines) is only rarely detected.

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