LANDAU-GINZBURG/CALABI-YAU CORRESPONDENCE,
GLOBAL MIRROR SYMMETRY AND ORLOV EQUIVALENCE

by ALESSANDRO CHIODO, HIROSHI IRTANI, and YONGBIN RUAN

ABSTRACT

We show that the Gromov-Witten theory of Calabi-Yau hypersurfaces matches, in genus zero and after an analytic continuation, the quantum singularity theory (FJRW theory) recently introduced by Fan, Jarvis and Ruan following a proposal of Witten. Moreover, on both sides, we highlight two remarkable integral local systems arising from the common formalism of $\hat{\Gamma}$-integral structures applied to the derived category of the hypersurface $\{W = 0\}$ and to the category of graded matrix factorizations of $W$. In this setup, we prove that the analytic continuation matches Orlov equivalence between the two above categories.

CONTENTS

1. Introduction ................................................................. 128
   1.1. Overview .......................................................... 128
   1.2. Mirror symmetry .................................................. 132
   1.3. Plan of the paper ................................................ 133
2. $\hat{\Gamma}$-integral structure and main statements ................................ 133
   2.1. FJRW theory ...................................................... 133
   2.2. GW theory ....................................................... 140
   2.3. Quantum cohomology and quantum connection ......................... 142
   2.4. Quantum D-modules and integral structure .......................... 147
   2.5. Statements of the main results ................................... 153
   2.6. Flowchart of the proof .......................................... 157
3. Computing FJRW theory .................................................. 157
   3.1. Extending FJRW theory ......................................... 158
   3.2. Twisted FJRW theory and Givental's formalism ...................... 160
   3.3. Family of elements on the Lagrangian cone ........................ 163
   3.4. The twist by the equivariant Euler class .......................... 164
   3.5. The $\epsilon_G$-twisted quantum connection ........................ 167
4. Orlov equivalence matches Mellin-Barnes analytic continuation ............... 171
   4.1. Matrix factorizations .......................................... 171
   4.2. Orlov equivalence .............................................. 174
   4.3. Twisted I-functions and Mellin-Barnes continuation ............... 177
   4.4. The non-equivariant limit and Orlov equivalence ................... 183
5. Construction of global D-module ......................................... 186
   5.1. Multi-GKZ system ............................................... 186
   5.2. Refined mirror theorem ........................................ 189
   5.3. Analytic continuation $U^w$ revisited ............................ 196
   5.4. The non-equivariant limit and its reduction ....................... 198
   5.5. Reconstruction of the big quantum D-module ....................... 203
   5.6. Monodromy and autoequivalences ................................ 205
Acknowledgements ................................................................ 208
Appendix A: Proof of Proposition 2.1 ......................................... 208
Appendix B: Compatibility with FJRW setup .................................. 212
References ................................................................. 214

DOI 10.1007/s10240-013-0056-z
1. Introduction

1.1. Overview. — The so-called Landau-Ginzburg/Calabi-Yau correspondence (LG/CY correspondence for short) in string theory [30, 51, 67] describes a relationship between the sigma model on a Calabi-Yau hypersurface and the Landau-Ginzburg model whose potential is the defining equation of the Calabi-Yau. In Witten’s gauged linear sigma model [69], the LG/CY correspondence arises, roughly speaking, from a variation of GIT quotient.

Let $w_1, \ldots, w_N$ be coprime positive integers and $x_1, \ldots, x_N$ be variables of degree $w_1, \ldots, w_N$. Let $W(x_1, \ldots, x_N)$ be a weighted homogeneous polynomial of degree $d$ which has an isolated critical point only at the origin. We assume (i) the Calabi-Yau condition $d = w_1 + \cdots + w_N$ and (ii) the Gorenstein condition: $w_j$ divides $d$ for all $1 \leq j \leq N$.

In this paper, we discuss two objects:

- The Calabi-Yau hypersurface $X_W = \{ W = 0 \}$ in the weighted projective space $\mathbb{P}(w) = \mathbb{P}(w_1, \ldots, w_N)$. This is quasi-smooth (i.e. smooth in the sense of stacks) by the assumption on $W$ above.

- The Landau-Ginzburg orbifold $(\mathbb{C}^N, W, \mu_d)$. It consists of the space $\mathbb{C}^N$ equipped with an action of $\mu_d = \{ g \in \mathbb{C}^\times \mid g^d = 1 \}$, $(x_1, \ldots, x_N) \mapsto (g^{w_1} x_1, \ldots, g^{w_N} x_N)$ and a $\mu_d$-invariant function $W: \mathbb{C}^N \to \mathbb{C}$.

These two models arise from the following GIT quotient. Consider the $\mathbb{C}^\times$-action on the vector space $\mathbb{C}^N \times \mathbb{C}$ with co-ordinates $(x_1, \ldots, x_N, p)$:

$$(x_1, \ldots, x_N, p) \mapsto (t^{w_1} x_1, \ldots, t^{w_N} x_N, t^{-d} p), \quad t \in \mathbb{C}^\times.$$  

We endow the space $\mathbb{C}^N \times \mathbb{C}$ with the $\mathbb{C}^\times$-invariant potential $\tilde{W}(x, p) := p W(x)$. There are two possible GIT quotients of this space: one is the quotient of $(\mathbb{C}^N \setminus \{0\}) \times \mathbb{C}$ and

\[ W(x_1, \ldots, x_N) \text{ weighted homogeneous polynomial (Section 1.1)} \\
X_W \text{ Calabi-Yau hypersurface defined by } W \text{ in } \mathbb{P}(w) \text{ (Section 1.1)} \\
\mu_d \text{ the group of } d\text{-th roots of unity} \\
H \text{ state space } (H(W, \mu_d) \text{ or } H_{CR}(X_W), \text{Sections 2.1.1, 2.2.1}) \\
H' \text{ narrow/ambient part (Section 2.3.1)} \\
H'' \text{ broad/primitive part (Section 2.3.1)} \\
\{t^p\} \text{ linear co-ordinates on the state space associated to a basis } \{T_i\} \text{ (Sections 2.3, 2.4, 3.5.2)} \\
\hat{\Gamma} \text{ Gamma class (Section 2.4.4)} \\
\overline{H} \text{ state space of twisted theories } (H_{ext} \text{ or } H_{CR}(\mathbb{P}(w)), \text{Section 3.5.2}) \\
M \text{ global Kähler moduli space } (\mathbb{P}(1, d) \setminus \{0, v_c\}, \text{Section 5.1})
the other is the quotient of $\mathbb{C}^N \times (\mathbb{C} \setminus \{0\})$. In the former case, we get the total space of the line bundle $\mathcal{O}(-d) \to \mathbb{P}(w)$ endowed with the function $\widetilde{W}$. This should reduce to the sigma model on the Calabi-Yau hypersurface $X_W$. In the latter case, we get the Landau-Ginzburg orbifold $(\mathbb{C}^N, W, \mu_d)$.

The GIT quotient itself does not change inside a “chamber” of stability parameters, but the actual physical theory depends on a continuous and complexified stability parameter $r + i\theta \in \mathbb{C}$. The CY theory arises for $r \to \infty$ and the LG theory for $r \to -\infty$. The stability parameter $r + i\theta$ varies along a complex manifold $\mathcal{M}$ called the global Kähler moduli space. In the case at hand, it is identified with (a Zariski open subset of) the weighted projective line $\mathbb{P}(1, d)$. The local picture near the $\mu_d$-point in $\mathbb{P}(1, d)$ corresponds to the LG model above and the $\mu_d$-point is called LG point. The local picture near the antipodal point corresponds to the CY geometry and the antipodal point is called large radius limit point. These points are interesting asymptotically: we often work on punctured discs centered on them and refer to them as limit points (see Figure 1). There is another limit point in the Kähler moduli space where the mirror has a conifold singularity; by abuse of terminology, we refer to this limit point as a conifold point.

This paper is concerned with two aspects of topological string theory: the category of D-branes of type B (B-branes) and the closed string theory of type A (A-model). In this paper, the term “B-brane” (or “brane”) has a precise mathematical meaning. On the CY side, the category of B-branes is the derived category of coherent sheaves on the Calabi-Yau hypersurface $X_W$. On the LG side, the category of B-branes is identified with the category $\text{MF}^\text{gr}_{\mu_d}(W)$ of graded matrix factorizations of $W$ [37, 52, 68]. On the other hand, the mathematical A-model on the Calabi-Yau $X_W$ is given by GW (Gromov-Witten) theory. The mathematical A-model for the Landau-Ginzburg orbifold was formulated recently by Fan-Jarvis-Ruan [26] as the intersection theory on the moduli space of W-spin curves. This is called FJRW (Fan-Jarvis-Ruan-Witten) theory. About these theories, the following LG/CY correspondences are known in mathematics:

1. **B-brane LG/CY correspondence**: Orlov [53] constructed derived equivalences $\Phi_l$ between the categories of B-branes indexed by an integer $l \in \mathbb{Z}$:

   $$\Phi_l : \text{D}^b(X_W) \cong \text{MF}^\text{gr}_{\mu_d}(W).$$
The purpose of this paper is to extend the correspondence (2) to a general weighted homogeneous polynomial $W$ and to describe a relationship between (1) and (2).

More precisely, “analytic continuation” in (2) means the following. In genus zero, GW theory and FJRW theory yield quantum D-modules over small neighbourhoods of the corresponding limit points; we show that these are restrictions of a certain global D-module over the Kähler moduli space $\mathcal{M}$. Note that the category of B-branes should be independent of the Kähler structure on $X_W$. Hence B-branes are “locally constant” data over the Kähler moduli space around the limit point in each theory. In fact, we associate to each B-brane a flat section of the quantum D-module. The flat section here is asymptotic (in the limit point) to the Chern character of the brane multiplied by the $\hat{\gamma}/\Gamma_1$-class. This defines a $\mathbb{Z}$-local system underlying the quantum D-module whose fibre is the numerical K-group of the category of B-branes. We call it the $\hat{\gamma}/\Gamma_1$-integral structure of the quantum D-module. This has been introduced for GW theory by Iritani [42] and Katzarkov-Kontsevich-Pantev [44]. Here the role of the $\hat{\gamma}$-class (see Definition 2.17) is to preserve the Euler pairing $\chi(E, F) := \sum_{i \in \mathbb{Z}} \dim \text{Hom}(E, F[i])$ in B-brane categories. Indeed, $\hat{\gamma}$ can be regarded as a square root of the Todd class. When $X_W$ is a manifold, we have

$$((-1)^{\deg \hat{\gamma}_{X_W}}) \cdot \hat{\gamma}_{X_W} = (2\pi i)^{\deg \text{Td}_{X_W}}$$

thanks to the functional equation $\Gamma(1-z)\Gamma(1+z) = \pi z / \sin(\pi z)$. Thus one can think of our flat section associated to a B-brane as a quantum version of the Mukai vector. In this paper, we extend the $\hat{\gamma}$-integral structure to FJRW theory. Our main results are stated as follows. We refer the reader to Theorems 2.23, 2.25 for more precise statements.

**Theorem 1.1.**

(i) The ambient part quantum D-module of $X_W$ and the narrow part quantum D-module of $(\mathbb{C}^N, W, \mu_d)$ are analytically continued to each other, i.e., both of them are the restrictions of a global D-module over the Kähler moduli space $\mathcal{M}$.

(ii) The analytic continuation in (i) matches up the $\hat{\gamma}$-integral structures on both quantum D-modules. Moreover, the induced isomorphism between the numerical K-groups of the categories of B-branes coincide with the one induced by the Orlov derived equivalence.

---

2 In fact, it coincides with the Mukai vector for K3 surfaces.

3 See Section 2.4.2 for the definition of the quantum D-modules. We mean by “ambient part” the cohomology classes pulled back from the ambient space, see Section 2.2.1; in FJRW side, this has a counterpart called “narrow part”, see Section 2.1.1.
A prototype of our result is the work of Borisov-Horja [5], where they showed that the analytic continuation of the GKZ hypergeometric system is induced from a Fourier-Mukai transformation between the $K$-groups of toric Calabi-Yau orbifolds, under a suitable identification of the spaces of local solutions with the $K$-groups. In our case, the GKZ system is replaced with the quantum D-modules of GW/FJRW theories and the Fourier-Mukai transformation is replaced with the Orlov equivalence.

Since the global D-module over $\mathcal{M}$ have non-trivial monodromies, the analytic continuation of flat sections depends on the choice of (a homotopy type of) a path. On the other hand, the Orlov equivalence $\Phi_l$ depends on an integer $l \in \mathbb{Z}$. The recent physics paper [34] by Herbst-Hori-Page clarified (by a physical argument) the dependence of a derived equivalence on the choice of a path. We confirm the prediction of [34] that a path $\gamma$ passing through the $l$th “window” corresponds to the $l$th Orlov equivalence $\Phi_l$.

Moreover, we check that the monodromy representation of the fundamental group of $\mathcal{M} = \mathbb{P}(1,d) \setminus \{2 \text{ points}\}$ factors through the group of autoequivalences of $D^b(\mathbb{X}_W)$. The following theorem refines Part (ii) of Theorem 1.1.

\textbf{Theorem 1.2} (Theorem 2.26).

(i) For each integer $l \in \mathbb{Z}$, there exists a path $\gamma_l$ from a neighbourhood of the large radius limit point to a neighbourhood of the LG point such that the analytic continuation along $\gamma_l^{-1}$ is induced by the Orlov equivalence $\Phi_l$.

(ii) Let $N'(X_W)$ be the subgroup (23) of the numerical $K$-group of $X_W$ consisting of $K$-classes whose Chern characters lie in the ambient cohomology $H_{\text{amb}}(X_W)$ and let $\chi$ be the Euler pairing. The monodromy representation of the global D-module in Theorem 1.1

$$\rho : \pi_1(\mathcal{M}) \to \text{Aut}(N'(X_W), \chi)$$

can be lifted to a group homomorphism

$$\hat{\rho} : \pi_1(\mathcal{M}) \to \text{Auteq}(D^b(X_W))/[2],$$

where $[2]$ is the 2-shift functor. The homomorphism $\hat{\rho}$ sends a (clockwise) loop around the conifold point to the spherical twist by the structure sheaf.

Since the work of Seidel-Thomas [63], the monodromy group action on $D^b(X)$ has been widely studied. Horja [38] identified the conifold monodromy of the GKZ system with the spherical twist. Aspinwall [3, Section 7.1.4] observed that the 5th power of the monodromy around the LG point corresponds to the 2-shift (for a quintic). We deduce the existence of the lift $\hat{\rho}$ from a result of Canonaco-Karp [9]. The above theorem suggests an autoequivalence group action on GW theory. However we do not know if $\hat{\rho}$ is injective. The induced homomorphism $\rho$ is never injective when $\dim X_W$ is even (since the conifold monodromy is involutive), but it is still possible that $\rho$ is injective for an odd-dimensional Calabi-Yau $X_W$. 


1.2. Mirror symmetry. — The interaction between B-branes and the A-model above can be explained most clearly via mirror symmetry. Here we consider Hodge-theoretic mirror symmetry, Kontsevich’s homological mirror symmetry [46] and their mutual relationships. See also [15] for the discussion on global mirror symmetry for finite group quotients of Calabi-Yau hypersurfaces.

The mirror of $X_W$ is given by a certain Calabi-Yau compactification $Y_v$ (Batyrev’s mirror [4]) of the affine variety

$$Y_v^\circ := \{ (x_1, \ldots, x_N) \in (\mathbb{C}^*)^N \mid x_1 + \cdots + x_N = 1, x_1^{w_1} x_2^{w_2} \cdots x_N^{w_N} = v \}$$

where the parameter $v$ is identified with an inhomogeneous co-ordinate of $\mathcal{M} = \mathbb{P}(1, d) \setminus \{ 2 \text{ points} \}$ such that $v = 0$ is the large radius limit and that $v = \infty$ is the LG point. The mirror $Y_v$ may have Gorenstein terminal quotient singularities. Note that $\mathcal{M}$ now plays a role of the complex moduli of $Y_v$. Under mirror symmetry, the category of B-branes should be equivalent to the category of A-branes of the mirror. Mathematically, the category of A-branes is the derived Fukaya category whose objects are (twisted complexes of) graded Lagrangian submanifolds. Likewise, the A-model theory should be equivalent the B-model theory of the mirror, which is, at genus zero, the variation of Hodge structure associated to the deformation of the complex structure. We get the mirror statements of (1) and (2).

(1') A-brane “mirror LG/CY” correspondence: The derived Fukaya category of $Y_v$ is independent of $v \in \mathcal{M}$.

(2') B-model “mirror LG/CY” correspondence: There exists a global variation of Hodge structure (VHS) $H^{N-2}(Y_v) = \bigoplus_{p+q=N-2} H^{p,q}(Y_v)$ over $\mathcal{M}$.

Because $Y_v$ does not change as a symplectic manifold (or orbifold) as $v$ varies, the Fukaya category should be independent of $v$ (if it is defined). The B-model VHS is tautologically “analytically continued” over $\mathcal{M}$. Moreover, the category of A-branes and the B-model have a natural integration pairing. Namely, one can integrate a de Rham cohomology class on $Y_v$ over a Lagrangian submanifold. By this pairing, an A-brane (a Lagrangian submanifold) gives rise to a dual flat section of the B-model VHS, i.e. a middle homology class in $H_{N-2}(Y_v, \mathbb{Z})$ represented by the brane. This is exactly dual to the phenomenon we described in Section 1.1.

The $\hat{\Gamma}$-integral structure in GW theory for $X_W$ is actually mirrored from the natural integral structure in the B-model of $Y_v$ (see also [15, Conjecture 4.2.10]).

Theorem 1.3 ([43, Theorem 6.9]). — The ambient A-model VHS of a Calabi-Yau hypersurface $X_W$ equipped with the ambient $\hat{\Gamma}$-integral structure is isomorphic to the residual B-model VHS of $Y_v$ equipped with the vanishing cycle integral structure near $v = 0$ under the mirror map $\tau_{GW}: \{ |v| < \epsilon \} \rightarrow H^2_{\text{amb}}(X_W)/(G)$ in Theorem 2.23.

Here, the ambient A-model VHS is the ambient part quantum D-module in Theorem 1.1 restricted to $z = 1$ (see Remark 2.13); the ambient $\hat{\Gamma}$-integral structure is the
Z-local system consisting of flat sections associated to vector bundles on X_W which are restricted from the ambient space P(w). The residual B-model VHS is defined to be the pure part Gr_W^{N-2} H^{N-2}(Y_v) \subset H^{N-2}(Y_v) of the Deligne mixed Hodge structure of the affine variety Y_v, the vanishing cycle integral structure on it is spanned by the Poincaré duals of vanishing cycles of the function x_1 + \cdots + x_N on the torus \{(x_1, \ldots, x_N) \in (\mathbb{C}^*)^N \mid \prod_{i=1}^{N} x_i^w = v\}. See [43] for the details. Because K-classes of vector bundles restricted from P(w) correspond, under Orlov equivalence, to the K-classes of graded Koszul matrix factorizations (Proposition 4.11), we have the following corollary (see also [15, Conjecture 4.2.11]).

**Corollary 1.4.** — The narrow A-model VHS of the Landau-Ginzburg model (\mathbb{C}^N, W, \mu_d) equipped with the subsystem of the \hat{\Gamma}-integral structure spanned by K-classes of graded Koszul matrix factorizations is isomorphic to the residual B-model VHS of Y_v equipped with the vanishing cycle integral structure near v = \infty under the mirror map \tau_{FJRW} : \{|v|^{-1/d} < \epsilon\} \to H^2_{nar}(W, \mu_d)/\langle G \rangle in Theorem 2.23.

In particular, both the quantum D-module of X_W and of (\mathbb{C}^N, W, \mu_d) over the image of the mirror map give a polarized variation of Z-Hodge structure.

1.3. Plan of the paper. — In Section 2, we introduce the \hat{\Gamma}-integral structure on the quantum D-module associated to FJRW and GW theories. Then we state our main theorems in a precise way. In Section 3, we introduce twisted FJRW invariants and calculate the (twisted) I-function of FJRW theory. This provides the main ingredients of the paper. In Section 4, we calculate the analytic continuation of the I-function and show that the connection matrix matches the Orlov equivalence. In Section 5, we construct a global D-module over the Kähler moduli and prove the main theorems.

2. \hat{\Gamma}-integral structure and main statements

We briefly review FJRW (Fan-Jarvis-Ruan-Witten) theory for (W, \mu_d) and GW (Gromov-Witten) theory for X_W and introduce the \hat{\Gamma}-integral structure on the quantum D-modules of both theories. Then we give a precise statement of the main results.

2.1. FJRW theory. — The FJRW invariants “count” the number of solutions to a non-linear PDE, the so-called Witten equation. These define a cohomological field theory on the FJRW state space. In this paper, we restrict ourselves to the genus zero FJRW invariants from the “narrow part”. In this case, the Witten equation has only trivial solutions and the invariants reduce to intersection numbers of tautological classes on the moduli space of d-spin curves. For the details of the full FJRW theory, we refer the reader to the original articles [25, 26].
2.1.1. **State space.** — Let \((\mathbf{C}^N, W, \mu_d)\) be the Landau-Ginzburg orbifold in the previous Section 1.1. Let \(\zeta := \exp(2\pi i/d) \in \mu_d\) denote a primitive \(d\)th root of unity. Let \((\mathbf{C}^N)_k\) denote the \(\zeta^k\)-fixed subspace of \(\mathbf{C}^N\) and \(W_k : (\mathbf{C}^N)_k \to \mathbf{C}\) denote the restriction of \(W\). We also write \(N_k = \dim_{\mathbf{C}}(\mathbf{C}^N)_k\). The FJRW state space \(\mathcal{H}(W, \mu_d)\) is defined to be
\[
\mathcal{H}(W, \mu_d) := \bigoplus_{k=0}^{d-1} \mathcal{H}(W, \mu_d)_k
\]
where the sector \(\mathcal{H}(W, \mu_d)_k\) associated to \(\zeta^k \in \mu_d\) is given by
\[
\mathcal{H}(W, \mu_d)_k := \mathcal{H}_{\mathcal{N}}((\mathbf{C}^N)_k, W_k^{\pm \infty}, \mathbf{C})^\mu_d,
\]
\[
W_k^{\pm \infty} := \{ x \in (\mathbf{C}^N)_k : \pm \Im(W_k(x)) > 0 \}.
\]
The degree of an element \(\phi \in \mathcal{H}(W, \mu_d)_k\) is defined to be
\[
\deg \phi := N_k + 2 \sum_{i=1}^{N} \langle kq_i \rangle - 2
\]
where \(q_i := w_i/d\). This is an integer since \(\sum_{i=1}^{N} q_i = 1\). Let \(\langle \cdot, \cdot \rangle\) denote the natural intersection pairing
\[
\langle \cdot, \cdot \rangle : \mathcal{H}_{\mathcal{N}}((\mathbf{C}^N)_k, W_k^{\pm \infty}; \mathbf{C}) \times \mathcal{H}_{\mathcal{N}}((\mathbf{C}^N)_k, W_k^{\pm \infty}; \mathbf{C}) \to \mathbf{C}
\]
and \(I : \mathbf{C}^N \to \mathbf{C}^N\) denote the map \((x_1, \ldots, x_N) \mapsto (\bar{\zeta}^{w_1} x_1, \ldots, \bar{\zeta}^{w_N} x_N)\) for \(\bar{\zeta} = \exp(\pi i/d)\).

Because \(W(I(x)) = -W(x)\), we have a map
\[
\Gamma^* : \mathcal{H}(W, \mu_d)_{d-k} \cong \mathcal{H}_{\mathcal{N}}((\mathbf{C}^N)_k, W_k^{\pm \infty}; \mathbf{C})^\mu_d \to \mathcal{H}_{\mathcal{N}}((\mathbf{C}^N)_k, W_k^{\pm \infty}; \mathbf{C})^\mu_d.
\]
We define the pairing between \(\alpha_1 \in \mathcal{H}(W, \mu_d)_k\) and \(\alpha_2 \in \mathcal{H}(W, \mu_d)_{d-k}\) by
\[
\langle \alpha_1, \alpha_2 \rangle := \frac{1}{d} \langle \alpha_1, \Gamma^* \alpha_2 \rangle.
\]
Setting \(\langle \alpha_1, \alpha_2 \rangle = 0\) for \(\alpha_1 \in \mathcal{H}(W, \mu_d)_k\), \(\alpha_2 \in \mathcal{H}(W, \mu_d)_l\), with \(k + l \neq d\), we obtain a graded symmetric non-degenerate pairing \((\cdot, \cdot)\) on the state space \(\mathcal{H}(W, \mu_d)\). The pairing in this paper differs from that in [26] by the factor \(1/d = 1/|\mu_d|\). See Appendix B for this convention.

We say that a sector \(\mathcal{H}(W, \mu_d)_k\) is **narrow** if \((\mathbf{C}^N)_k = \{0\}\) and **broad** otherwise.

Each narrow sector \(\mathcal{H}(W, \mu_d)_k\) is one-dimensional and we denote by \(\phi_{k-1} \in \mathcal{H}(W, \mu_d)_k\) the standard basis given as the identity class on \((\mathbf{C}^N)_k = \{0\}\). We set
\[
\text{Nar} := \{ 0 \leq k \leq d-1 : (\mathbf{C}^N)_k = \{0\} \text{ i.e. } kq_i \notin \mathbf{Z} \text{ for all } i \}
\]

---

4 Fan-Jarvis-Ruan originally called these sectors “Neveu-Schwarz” and “Ramond” respectively, but they changed the names later.

5 Note the shift of the index \(k\) by one. An element \(\phi_k\) with \(k+1 \notin \text{Nar}\) will be introduced later in Section 3 in the context of “extended theory”, but notice that \(\phi_k\) with \(k+1 \notin \text{Nar}\) is not an element of the original FJRW state space.
and define the narrow part as
\[ H_{\text{narrow}}(W, \mu_d) = \bigoplus_{k \in \text{Nar}} H(W, \mu_d)_k = \bigoplus_{k \in \text{Nar}} C_\phi_{k-1}. \]

Note that \( \deg \phi_k = 2 \sum_{j=1}^N \langle kq_j \rangle \) for \( k + 1 \in \text{Nar} \). The degree zero element \( \phi_0 \in H(W, \mu_d)_1 \) plays the role of the identity in the FJRW theory. The pairing \((\cdot, \cdot)\) restricts to a non-degenerate pairing on \( H_{\text{narrow}}(W, \mu_d) \)

\[ (\phi_k, \phi_l) = \frac{1}{d} \delta_{d,(k+1)+l+1}, \quad k + 1, l + 1 \in \text{Nar} \]

and \( H_{\text{narrow}}(W, \mu_d) \) is orthogonal to the broad part \( H_{\text{broad}}(W, \mu_d) := \bigoplus_{k \notin \text{Nar}} H(W, \mu_d)_k \).

For a polynomial \( f \) on \( C^n \), we define the Jacobi space \( \Omega(f) \) by

\[ \Omega(f) := \Omega^n_{C^n}/df \wedge \Omega^{n-1}_{C^n}, \]

where \( \Omega^k_{C^n} \) denotes the space of algebraic \( k \)-forms on \( C^n \). When \( f \) has an isolated critical point at the origin, we have the Grothendieck residue pairing \( \text{Res}_f : \Omega(f) \otimes \Omega(f) \to C \) (see [31]):

\[ \text{Res}_f([a(y)dy], [b(y)dy]) := \text{Res} \left[ \frac{a(y)b(y)dy}{\partial_1 f, \ldots, \partial_n f} \right], \]

where \( y = (y_1, \ldots, y_n) \) is a co-ordinate system on \( C^n \) and \( dy = dy_1 \wedge \cdots \wedge dy_n \). The residue pairing is independent of the choice of co-ordinates. We shall consider the Jacobi space \( \Omega(W) \) associated to the homogeneous polynomial \( W \) on \( (C^N)_k \). The space \( \Omega(W) \) is graded by the usual degree: \( \deg x_j = \deg dx_j = w_j \).

**Proposition 2.1.** — We have a canonical isomorphism

\[ H(W, \mu_d)_k \cong \Omega(W)_k^{\mu_d}. \]

Under this isomorphism, the pairing \((\cdot, \cdot) : H(W, \mu_d)_k \times H(W, \mu_d)_{d-k} \to C \) translates into the Grothendieck residue pairing between \( \Omega(W)_k^{\mu_d} \) and \( \Omega(W)_{d-k}^{\mu_d} \cong \Omega(W)_k^{\mu_d} \):

\[ [\varphi] \otimes [\psi] \mapsto (-1)^{\frac{N(N-1)}{2}} (2\pi i)^N \frac{1}{d} \text{Res}_{W_k}([\varphi], (-1)^{[\psi]}[\psi]), \]

where \( [\psi] \) is the degree of \( \psi \in \Omega(W)_{d-k}^{\mu_d} \) divided by \( d \). Notice that the above isomorphism does not match up the grading on the FJRW state space \( H(W, \mu_d)_k \) (which is homogeneous of degree \( N_k + 2(\sum_{j=1}^N (kq_j) - 1) \)) and the grading on the Jacobi space \( \Omega(W)_k^{\mu_d} \).

The isomorphism (7) is given by the Hodge decomposition (95). See Appendix A for the proof. This description is used in Section 4.1.1 (and in Section 2.4.4) to discuss the Chern character and Riemann-Roch for matrix factorizations.

\[ ^6 \text{It is isomorphic to the Jacobi ring } C[y_1, \ldots, y_n]/(\partial f, \ldots, \partial_d f), \text{ but notice that } \Omega(f) \text{ is not a ring.} \]
2.1.2. Moduli of $W$-curves. — In this paper all stacks are defined over $\mathbb{C}$. By a pointed orbicurve, we mean a proper and connected one-dimensional Deligne-Mumford stack $C$ which has only nodes as singularities, which is equipped with distinct marked points $\sigma_1, \ldots, \sigma_n$ on the smooth locus, and which has possibly non-trivial stabilizers only at the marked points and the nodes. For a positive integer $d$, a pointed orbicurve $C$ is called $d$-stable [11] if the associated pointed coarse curve $|C|$ is stable and if all the stabilizers at the nodes and the marked points are isomorphic to $\mu_d$. We always assume that every node of an orbicurve is balanced [2], i.e. formally locally near a node, the curve is isomorphic to $\left[ \text{Spec} \left( \mathbb{C}[x,y]/xy \right) / \mu_d \right]$, where the action of $\mu_d$ is given by $(x,y) \mapsto (\zeta x, \zeta^{-1} y)$. For a pointed orbicurve $(C, \sigma_1, \ldots, \sigma_n)$, we introduce an invertible sheaf $\omega_{\text{log}}$ on $C$ which may be regarded as the dualizing sheaf $\omega_C$ twisted at the stacky markings $\sigma_1, \ldots, \sigma_n$ (each of which is of degree $1/d$) or equivalently as the pullback of the dualizing sheaf $\omega_{|C|}$ twisted at the coarse markings $|\sigma_1|, \ldots, |\sigma_n|$ (each of which is of degree one)

$$\omega_{\text{log}} = \omega_C(\sigma_1 + \cdots + \sigma_n) = q^*\left( \omega_{|C|}(|\sigma_1| + \cdots + |\sigma_n|) \right)$$

where $q: C \to |C|$ is the natural map. In other words, $\omega_{\text{log}}$ is the sheaf of logarithmic differential forms on $C$ with poles only at marked points and nodes, and such that the sum of the residues at each node is zero.

A $d$-spin structure on a pointed orbicurve $C$ is a line bundle $L \to C$ together with an isomorphism $\varphi: L^\otimes d \cong \omega_{\text{log}}$. Write $W(x_1, \ldots, x_N) = \sum_{i=1}^l c_i \prod_{j=1}^N x_j^{m_{ij}}$, where $\prod_{j=1}^N x_j^{m_{ij}}$, $i = 1, \ldots, l$ are mutually distinct monomials and $c_i \neq 0$. A $W$-structure on a pointed orbicurve $C$ is a collection of line bundles $L_1, \ldots, L_N$ (corresponding to the variables $x_1, \ldots, x_N$) on $C$ together with isomorphisms

$$\varphi_i: \bigotimes_{j=1}^N L_j^\otimes m_{ij} \cong \omega_{\text{log}}, \quad i = 1, \ldots, l.$$

This generalizes the notion of a $d$-spin structure (see Remark 2.2). Since $W$ is weighted homogeneous of degree $d$, a $d$-spin structure $L \to C$ gives rise to a $W$-structure by setting

$$L_i = L^\otimes w_i, \quad i = 1, \ldots, N.$$ 

A $W$-structure does not necessarily arise from a $d$-spin structure in this way. In this paper, however, we restrict our attention to a $W$-structure coming from a $d$-spin structure.\footnote{This is because we are interested in a generic weighted homogeneous polynomial $W$. More precisely, we can add to $W$ a weighted homogeneous Laurent polynomial $Z$ so that the group of diagonal symmetries preserving $W + Z$ is exactly $\mu_d$; then every $(W + Z)$-structure comes from a $d$-spin structure. This means that the group $\mu_d$ is admissible in the sense of \cite[Section 2.3]{26}.}

Let $L$ be a $d$-spin structure on a $d$-stable pointed orbicurve $C$. The stabilizer $\mu_d$ at a marked point $\sigma$ acts on the fibre $L_\sigma$ via a homomorphism $\mu_d \to \mathbb{C}^\times$, which is of the form $t \mapsto t^k$ for a unique $0 \leq k < d$. We call the rational number $\text{age}_\sigma(L) := k/d \in \mathbb{Q}$.\footnote{This is because we are interested in a generic weighted homogeneous polynomial $W$. More precisely, we can add to $W$ a weighted homogeneous Laurent polynomial $Z$ so that the group of diagonal symmetries preserving $W + Z$ is exactly $\mu_d$; then every $(W + Z)$-structure comes from a $d$-spin structure. This means that the group $\mu_d$ is admissible in the sense of \cite[Section 2.3]{26}.}
[0, 1) the age of \( L \) at \( \sigma \). The generator \( \zeta \in \mu_d \) acts on the fibre of the associated \( W \)-structure \( (L^{\otimes w_1}, \ldots, L^{\otimes w_N}) \) at \( \sigma \) by \( (\zeta^{w_1}, \ldots, \zeta^{w_N}) \); hence in this case we regard the marked point \( \sigma \) as corresponding to the sector \( H(W, \mu_d)_k \). For \( 0 \leq k_1, \ldots, k_n \leq d - 1 \), let \( \text{Spin}^d_{0,n}(k_1, \ldots, k_n) \) denote the moduli stack of \( d \)-stable orbicurves \( C \) of genus zero and with \( n \) marked points \( \sigma_1, \ldots, \sigma_n \) endowed with a \( d \)-spin structure \( L \to C \) such that \( \text{age}_{\sigma_i}(L) = ((k_i + 1)/d) \).

\[
\text{Spin}^d_{0,n}(k_1, \ldots, k_n) = \left\{ (C; \sigma_1, \ldots, \sigma_n; L; \varphi : L^{\otimes d} \cong \omega_{\log}) \mid \text{age}_{\sigma_i}(L) = \left( \frac{k_i + 1}{d} \right) \right\} / \text{isom.}
\]

See [13, Appendix] for a precise definition of \( \text{Spin}^d_{0,n} \) as a fibred category, where it is denoted by \( \mathcal{R}_d \). More precisely the substack denoted by \( \mathcal{R}_d(e^{2\pi i/\Theta_1}, \ldots, e^{2\pi i/\Theta_n}) \) in [13] corresponds to \( (9) \) as soon as \( \Theta_j = (k_j + 1)/d \); the present choice of notation allows a more straightforward formula for the FJRW invariants, see (12) and the following footnote 9.

The moduli stack \( \text{Spin}^d_{0,n}(k_1, \ldots, k_n) \) is smooth, proper and of Deligne-Mumford type [11]. It is non-empty if and only if \( \chi(L) = \deg L - \sum_{i=1}^n \text{age}_{\sigma_i}(L) - 1 = (n - 2)/d - \sum_{i=1}^n ((k_i + 1)/d) - 1 \) is an integer (see [1, 45, 65] for Riemann-Roch for orbicurves). When non-empty, it is of dimension \( n - 3 \). (It differs from the Deligne-Mumford space \( \mathcal{M}_{0,n} \) of stable curves only because of the stabilizers.) It is clear that the definition (9) extends verbatim to higher genera; we write \( \text{Spin}^d_{g,n}(k_1, \ldots, k_n) \) for the similar moduli space at genus \( g \). We use it only in Section 2.1.3 where we recall the general setup of [26].

**Remark 2.2.** — The definition (8) of a \( W \)-structure originates from the Witten equation, which is a system of PDE for sections \( s_i \in C^\infty(C, L_i) \), \( i = 1, \ldots, N \):

\[
\partial_j s_j + \partial_j W(s_1, \ldots, s_N) = 0.
\]

The equation makes sense under (8) and a suitable choice of a Hermitian metric on \( L_i \) (see [25]). When all the marked points correspond to the narrow sector, the zero sections are only possible solutions to the Witten equation [25, Theorem 3.3.8] and the FJRW invariant is the Euler class of the obstruction bundle over the moduli space of \( W \)-curves. We briefly review the general case.

**2.1.3. FJRW invariants.** — We review the virtual fundamental class defining the invariants of “full” FJRW theory following [26]. The general FJRW invariants are defined by an analytic moduli space of solutions to the Witten equation (10). We shall see in the next Section 2.1.4 that the narrow sector invariants have an algebro-geometric definition. We stress that the invariants introduced later in Section 3.1 are different from the invariants of the full theory; they are a natural and computable extension of the narrow sector invariants.
The virtual cycle\(^8\)

\[
\bigoplus_{0 \leq k_1, \ldots, k_n \leq d-1} \left[ \text{Spin}^d_{g,n}(k_1, \ldots, k_n) \right]^{\text{vir}}
\]

of the fully developed FJRW theory is defined in all genera and lies in

\[
\bigoplus_{0 \leq k_1, \ldots, k_n \leq d-1} \left( \text{H}_{2D(k_1, \ldots, k_n)} \left( \text{Spin}^d_{g,n}(k_1, \ldots, k_n); \mathbf{C} \right) \right)
\]

\[
\otimes \bigotimes_{i=1}^{n} \text{H}_{N_{k_i}} \left( \left( \mathbf{C}^N \right)_{k_i}, W_{k_i}^{+\infty}; \mathbf{C} \right)^{\mu_d}
\]

where \(D(k_1, \ldots, k_n) = (3g - 3 + n) + \sum_{j=1}^{N} \chi(L \otimes u_j)\) for a \(d\)-spin structure \(L\) from \(\text{Spin}^d_{g,n}(k_1, \ldots, k_n)\). Regard the relative homology group \(\text{H}_{N_{k_i}} \left( \left( \mathbf{C}^N \right)_{k_i}, W_{k_i}^{+\infty}; \mathbf{C} \right)^{\mu_d}\) above as the dual of the sector \(\text{H}(W, \mu_d)\) of the state space. Then the virtual cycle defines a linear operator (see [26, Equation (63)]):

\[
\text{H}(W, \mu_d)^{\otimes n} \xrightarrow{[\text{Spin}^d_{g,n}(k_1, \ldots, k_n)]^{\text{vir}}} \text{H}_{2D(k_1, \ldots, k_n)} \left( \text{Spin}^d_{g,n}(k_1, \ldots, k_n); \mathbf{C} \right)
\]

assigning to \(\bigotimes_{i=1}^{n} \alpha_i\) a nonzero cycle only if \(\alpha_i\) lies in \(\text{H}^{N_{k_i}} \left( \left( \mathbf{C}^N \right)_{k_i}, W_{k_i}^{+\infty}; \mathbf{C} \right)^{\mu_d}\). In FJRW theory, such a cycle is pushed forward via the natural forgetful morphism

\[
\text{st}: \text{Spin}^d_{g,n}(k_1, \ldots, k_n) \rightarrow \overline{\mathcal{M}}_{g,n}
\]

forgetting the data \(L\) and \(\varphi\) and passing to the coarse stable curve corresponding to \((C; \sigma_1, \ldots, \sigma_n)\). After Poincare duality and multiplication by a factor \(f_g = [G]^s / \deg(\text{st}) = d/g\), Fan, Jarvis and Ruan obtain a cohomological field theory. We show in Appendix B that the multiplication by \(f_g\) can be removed in all genera once we use the natural pairing (4) from Chen-Ruan cohomology. In genus zero, this simply amounts to the fact that our invariants are \(1/f_0 = 1/d\) times the FJRW genus-zero invariants of [26]. By making explicit the cohomological field theory of [26, Definition 4.2.1] after removing \(f_g\) and by applying the definition of the correlators of [26, Definition 4.2.6], we obtain

\[
\left\{ \tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n) \right\}_{\text{FJRW}}^{(g,n)} := \left[ \left[ \text{Spin}^d_{g,n}(k_1, \ldots, k_n) \right]^{\text{vir}} \cap \prod_{i=1}^{n} \alpha_i \right] \cap \prod_{i=1}^{n} \psi_i^{b_i}
\]

for \(\alpha_i \in \text{H}^{N_{k_i}} \left( \left( \mathbf{C}^N \right)_{k_i}, W_{k_i}^{+\infty}; \mathbf{C} \right)^{\mu_d}\) and \(b_i \geq 0\). The class \(\psi_i\) is the first Chern class of the line bundle on \(\text{Spin}^d_{g,n}(k_1, \ldots, k_n)\) whose fibre at a point \((C; \sigma_1, \ldots, \sigma_n; L; \varphi)\) is the cotangent space \(T_{\sigma_i}[C]\) of the coarse curve \([C]\) at \(\sigma_i\). Note that, since these classes are the pullback of the standard \(\psi\) classes from \(\overline{\mathcal{M}}_{g,n}\), we can directly integrate on \(\text{Spin}^d_{g,n}(k_1, \ldots, k_n)\) short-circuiting the pushforward st* of [26, Definition 4.2.1].

---

\(^8\) \([\text{Spin}^d_{g,n}(k_1, \ldots, k_n)]^{\text{vir}}\) is denoted by \([\overline{\mathcal{M}}_{g,n}\left(W, (k_1, \ldots, k_n)\right)]^{\text{vir}}\) in [26].
2.1.4. Narrow part FJRW invariants. — The definition, further simplifies in genus zero and when the entries are narrow states \( \phi_k \) with \( k + 1 \in \text{Nar} \). Let \( \pi : \mathcal{C} \to \text{Spin}_{0,a}(k_1, \ldots, k_n) \) be the universal orbicurve and \( \mathcal{L} \to \mathcal{C} \) be the universal \( d \)-spin structure. If all the entries are narrow states, then the following lemma shows that \( H^0(\mathcal{C}, \mathcal{L}^\otimes_{w_j}) \) vanishes all \( j \) (and for every curve \( \mathcal{C} \)); hence \(-R\pi_* \bigoplus_{j=1}^N \mathcal{L}^\otimes_{w_j} = R^1\pi_* \bigoplus_{j=1}^N \mathcal{L}^\otimes_{w_j}\) is a vector bundle which we refer to as the obstruction bundle. In these cases the virtual homology cycle is simply the Poincaré dual cycle of the top Chern class of the obstruction bundle [26, Theorem 4.1.8, Concavity].

**Lemma 2.3.** — Suppose that \( k_i + 1 \in \text{Nar} \) for all \( i \). Then \( H^0(\mathcal{C}, \mathcal{L}^\otimes_{w_j}) \) vanishes for all \( j = 1, \ldots, N \) for \( (\mathcal{C}; \sigma_1, \ldots, \sigma_n; L; \varphi) \in \text{Spin}_{0,a}^d(k_1, \ldots, k_n) \). In particular, the obstruction bundle \( \bigoplus_{j=1}^N R^1\pi_*(\mathcal{L}^\otimes_{w_j}) \) over \( \text{Spin}_{0,a}^d(k_1, \ldots, k_n) \) is locally free of rank \( 2-N + \sum_{i=1}^n \sum_{j=1}^N \langle k_i q_j \rangle \).

**Proof.** — Let \( p : \mathcal{C} \to \overline{\mathcal{C}} \) denote the map forgetting the stack-theoretic structures at all the markings (but not at the nodes). Then

\[
p^*p_*(\mathcal{L}^\otimes_{w_j}) = \mathcal{L}^\otimes_{w_j} \otimes \mathcal{O}_C \left( -\sum_{i=1}^n \text{age}_{\sigma_i}(\mathcal{L}^\otimes_{w_j}) \sigma_i \right)
\]

\[
= \mathcal{L}^\otimes_{w_j} \otimes \mathcal{O}_C \left( -\sum_{i=1}^n d\langle (k_i + 1)q_j \rangle \sigma_i \right).
\]

Here we regard \( \sigma_i \) as a stacky divisor of degree \( 1/d \). Hence

\[
\mathcal{L}' := (p^*p_*(\mathcal{L}^\otimes_{w_j}))^{\otimes(q_j^{-1})} = \omega_{\text{log}} \otimes \mathcal{O}_C \left( -\sum_{i=1}^n dq_j^{-1} \langle (k_i + 1)q_i \rangle \sigma_i \right).
\]

Notice that \( q_j^{-1} = d/w_j \) is an integer by the Gorenstein condition and we have \( q_j^{-1} \langle (k_i + 1)q_j \rangle \geq 1 \) since \( k_i + 1 \in \text{Nar} \). Hence \( \mathcal{L}' \) is a subsheaf of the pull-back of \( \omega_{|\mathcal{C}|} \). Note that \( H^0(|\mathcal{C}|, \omega_{|\mathcal{C}|}) = 0 \) since \( |\mathcal{C}| \) is of genus zero. Hence \( H^0(\mathcal{C}, L') = 0 \) and thus \( H^0(\mathcal{C}, p^*p_*(\mathcal{L}^\otimes_{w_j})) = 0 \). This implies \( H^0(\mathcal{C}, \mathcal{L}^\otimes_{w_j}) = 0 \). Finally, by Riemann-Roch [1, 45, 65], \( R^1\pi_*(\mathcal{L}^\otimes_{w_j}) \) is locally free of rank

\[
-\chi(\mathcal{L}^\otimes_{w_j}) = -1 - \deg \mathcal{L}^\otimes_{w_j} + \sum_{i=1}^n \text{age}_{\sigma_i}(\mathcal{L}^\otimes_{w_j})
\]

\[
= -1 - (n-2)q_j + \sum_{i=1}^n \langle (k_i + 1)q_i \rangle = -1 + 2q_j + \sum_{i=1}^n \langle q_i k_i \rangle.
\]

The conclusion follows. \( \square \)
Remark 2.4. — The above lemma usually fails without the Gorenstein condition. This is the main reason that we restrict ourselves to the Gorenstein case.

We can now specialize (11) and get the narrow part descendant FJRW invariants as

\[
\langle \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \rangle_{0, n}^{\text{FJRW}} = \int_{[\text{Spin}_{0, n}(k_1, \ldots, k_n)]} \prod_{i=1}^n \psi_i^{b_i} \cup \prod_{j=1}^N \epsilon_{\text{top}}(R^1 \pi_* (L^{\otimes w}))
\]

where \( \phi_{k_1}, \ldots, \phi_{k_n} \in H_{\text{nar}}(W, \mu_d) \) (i.e. \( k_i + 1 \in \text{Nar} \)) and \( b_1, \ldots, b_n \geq 0 \). We sometimes omit \( \tau_0 \) from the notation, e.g. writing \( \langle \phi_{k_1}, \ldots, \phi_{k_n} \rangle_{0, n}^{\text{FJRW}} \) for \( \langle \tau_0(\phi_{k_1}), \ldots, \tau_0(\phi_{k_n}) \rangle_{0, n}^{\text{FJRW}} \). The FJRW invariants satisfy the homogeneity ([26, Dimension Axiom in Section 4.1])

\[
\langle \tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n) \rangle_{0, n}^{\text{FJRW}} = 0 \quad \text{unless} \quad \sum_{i=1}^n \left( b_i + \frac{1}{2} \deg \alpha_i \right) = n + \hat{c} - 3,
\]

where \( \hat{c} := N - 2 \) is the dimension of the Calabi-Yau hypersurface \( X_W \). Again the invariants (12) differ from the original definition\(^9\) [13, 26] by the factor of \( 1/d \). See Appendix B.

Remark 2.5. — Polishchuk and Vaintrob [57] recently constructed a purely algebraic cohomological field theory of singularities based on matrix factorizations. They constructed a fundamental matrix factorization on the moduli space which plays the same role as the virtual fundamental class. The role of matrix factorizations in our paper is different from theirs, but it would be interesting to study the relationships.

2.2. GW theory. — GW theory for orbifolds has been developed by Chen-Ruan [10] in symplectic category and Abramovich-Graber-Vistoli [1] in algebraic category. We will work in the algebraic category. For the details, we again refer the reader to these original articles.

2.2.1. State space. — The state space of orbifold GW theory is given by the Chen-Ruan cohomology group of the orbifold. We explain the case of the Calabi-Yau hypersurface \( X_W \subset P(w) \). Set

\[
\mathfrak{F} := \{ 0 \leq f < 1 | f w_j \in \mathbb{Z} \text{ for some } 1 \leq j \leq N \}
\]

\[
= \{ 0 \leq f < 1 | fd \in \mathbb{Z}, fd \notin \text{Nar} \}.
\]

In the second line we used the Gorenstein condition (i.e. \( w_j \mid d \)). An element \( f \in \mathfrak{F} \) gives rise to the stabilizer \( \exp(2\pi i f) \) along the substack \( P(w)_f \subset P(w) \):

\[
P(w)_f := \left[ (((C^N)_{\mathfrak{F}} \setminus \{0\})/\mathbb{C}^*) \right]
\]

\(^9\) In [13, Section 3.1, (15)], where the case \( d = 5 \) was discussed, the invariant \( \langle \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \rangle_{0, n}^{\text{FJRW}} \) was defined to be

\[
5(\prod_{i=1}^n \psi_i^{b_i} \cup \epsilon_{\text{top}}(R^1 \pi_* (L^{\otimes w}))) \cap [R_3(e^{2\pi i (k_1+1)/5}, \ldots, e^{2\pi i (k_n+1)/5})]; \text{ see also [13, Section 2.3, (14)]}.\]
where recall that \((C^N)_c\) is the subspace of \(C^N\) fixed by \(\xi^k = \exp(2\pi i k/d)\). The inertia stacks \(\mathcal{P}(w)\), \(\mathcal{I}X_W\) are defined to be

\[
\mathcal{P}(w) = \bigsqcup_{f\in\mathfrak{F}} \mathcal{P}(w)_f, \quad \mathcal{I}X_W = \bigsqcup_{f\in\mathfrak{F}} (\mathcal{P}(w)_f \cap X_W).
\]

The Chen-Ruan cohomology \(H_{CR}(X_w)\) is defined to be the cohomology of the inertia stack:

\[
H_{CR}(X_w) := H(\mathcal{I}X_W; C) = \bigoplus_{f\in\mathfrak{F}} H(\mathcal{P}(w)_f \cap X_W; C).
\]

The degree of \(\alpha \in H^k(\mathcal{P}(w)_f \cap X_W)\), as an element of \(H_{CR}(X_w)\), is defined to be

\[
\deg \alpha = k + 2 \sum_{j=1}^N \langle f w_j \rangle.
\]

Note that this is an integer since \(\sum_{j=1}^N f w_j = fd \in \mathbb{Z}\). Let \(\text{inv}: \mathcal{P}(w)_f \cong \mathcal{P}(w)_{(1-f)}\) denote the natural isomorphism. For \(\alpha_1 \in H(\mathcal{P}(w)_f \cap X_W)\) and \(\alpha_2 \in H(\mathcal{P}(w)_{(1-f)} \cap X_W)\), we define

\[
(\alpha_1, \alpha_2) = \int_{\mathcal{P}(w)_f \cap X_W} \alpha_1 \cup \text{inv}^* \alpha_2.
\]

We set \((\alpha_1, \alpha_2) = 0\) if \(\alpha_1 \in H(\mathcal{P}(w)_{f_1} \cap X_W)\) and \(\alpha_2 \in H(\mathcal{P}(w)_{f_2} \cap X_W)\) and \(f_1 + f_2 \notin \mathbb{Z}\). Then \((\cdot, \cdot)\) defines a graded symmetric non-degenerate pairing on \(H_{CR}(X_w)\).

The ambient part \(H_{amb}(X_w)\) is defined to be the image of the restriction map

\[
H_{amb}(X_w) := \text{Im}(i^*: H_{CR}(\mathcal{P}(w)) = H(\mathcal{I}P(w)) \to H(\mathcal{I}X_W) = H_{CR}(X_w)).
\]

Let \(1_f \in H(\mathcal{P}(w)_f \cap X_W)\) denote the identity class on \(\mathcal{P}(w)_f \cap X_W\) and \(p = c_1(\mathcal{O}(1))\) denote the hyperplane class on \(\mathcal{P}(w)\). The ambient part is spanned, as a \(C\)-vector space, by \(p^i 1_f\), \(0 \leq i \leq \#\{1 \leq j \leq N \mid f w_j \in \mathbb{Z}\} - 1, f \in \mathfrak{F}\). The pairing \((\cdot, \cdot)\) restricts to a non-degenerate pairing on \(H_{amb}(X_w)\) and \(H_{amb}(X_w)\) is orthogonal to the complementary primitive part \(H_{pri}(X_w) := \text{Ker}(i_*: H_{CR}(X_w) \to H_{CR}(\mathcal{P}(w)))\).

\textbf{Remark 2.6 (Comparison of state spaces).} — The FJRW and GW state spaces can be identified, up to Tate twist, with the relative Chen-Ruan cohomology (Chiodo-Nagel [12]):

\[
H(W, \mu_d) = H_{CR}\left(\left[\frac{C^N}{\mu_d}, [W^{-1}(1)/\mu_d]\right]\right)
\]

\[
H_{CR}(X_W) = H_{CR}(\mathcal{O}_{\mathcal{P}(w)}(-d), \hat{W}^{-1}(1)).
\]

The first identification follows immediately from the definition. The second identification follows from the Thom isomorphism. Here \(\hat{W}: \mathcal{O}_{\mathcal{P}(w)}(-d) \to C\) is the function in
Section 1.1. Note that the pairs \((O_{P(\omega)}(-d), \tilde{W}^{-1}(1)), ([C^N/\mu_d], [W^{-1}(1)/\mu_d])\) are connected by a variation of GIT quotients. Chiodo-Ruan [14] showed that there exists a bigraded isomorphism

\[ H^{p,q}(W, \mu_d) \cong H^{p,q}_{\text{CR}}(X_W) \]

which preserves the pairings on both sides. In this paper, we will construct a graded isomorphism (preserving the pairing)

\[ H^{p,p}_{\text{narrow}}(W, \mu_d) \cong H^{p,p}_{\text{ambient}}(X_W) \]

which depends on a point of the Kähler moduli space. (Note that the narrow/ambient part has no \((p, q)\)-Hodge component with \(p \neq q\).) See Remark 2.24.

2.2.2. GW invariants. — For \(n \geq 0\) and \(\beta \in H_2(|X_W|, \mathbb{Z})\), let \((X_W)_{0,n,\beta}\) denote the moduli stack of genus zero, \(n\)-pointed, degree \(\beta\) stable maps to \(X_W\) (it is denoted by \(K_{0,n}(X_W, \beta)\) in [1]). This carries a virtual fundamental class \([((X_W)_{0,n,\beta})_{\text{vir}}] \in H_*(((X_W)_{0,n,\beta}; Q)\) of degree \(2(n + \hat{c} - 3)\) with \(\hat{c} := N - 2 = \dim X_W\). We have the evaluation map at the \(i\)th marked point

\[ \text{ev}_i : (X_W)_{0,n,\beta} \to \bar{T}X_W \]

where \(\bar{T}X_W\) denotes the rigidified cyclotomic inertia stack (see [1]). Because the underlying complex analytic spaces of \(\bar{T}X_W\) and \(TX_W\) are the same, we can define the pull-back \(\text{ev}_i^*: H_{\text{CR}}(X_W) \to H_*(((X_W)_{0,n,\beta}); \mathbb{Q})\). Let \(\psi_i\) be the first Chern class of the line bundle over \((X_W)_{0,n,\beta}\) whose fibre at a stable map is the cotangent space of the coarse curve at the \(i\)th marked point. The orbifold GW invariant is defined to be

\[ \langle \tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n) \rangle^{GW}_{0,n,\beta} := \int_{((X_W)_{0,n,\beta})_{\text{vir}}} \prod_{i=1}^n \text{ev}_i^*(\alpha_i) \psi_i^{b_i}. \] (15)

Here \(\alpha_1, \ldots, \alpha_n \in H_{\text{CR}}(X_W)\) and \(b_1, \ldots, b_n \geq 0\). Again we sometimes omit \(\tau_0\) from the notation. The orbifold GW invariants are also homogeneous. The invariant (15) vanishes unless \(\sum_{i=1}^n (b_i + \frac{1}{2} \deg \alpha_i) = n + \hat{c} - 3\).

2.3. Quantum cohomology and quantum connection. — We can associate the quantum cohomology rings to both of FJRW and GW theories. The quantum ring of FJRW theory is defined over \(\mathbb{C}\), whereas the quantum ring of GW theory is defined over the Novikov ring \(\Lambda := \mathbb{C}[\llbracket \text{Eff} \rrbracket]\). It is the completion of the group ring \(\mathbb{C}[\text{Eff}]\) of the semigroup \(\text{Eff} \subset H_2(|X_W|, \mathbb{Z})\) consisting of classes of effective curves. For \(\beta \in \text{Eff}\), we denote by \(Q^\beta\) the corresponding element in \(\Lambda\). In Section 2.3.2 below, we see how the divisor equa-
tion reduces the ground ring to $\mathbb{C}$ (by setting $Q^\beta = 1$) for GW theory. The construction here is standard and can be applied to any (genus zero) cohomological field theories with homogeneity. See [49].

In order to describe the quantum rings of both theories in a uniform way, we use the following notation. Let $K$ denote the ground ring. It is $\mathbb{C}$ for FJRW theory and $\Lambda$ for GW theory. Let $H$ denote the state space. It is $H(W, \mu_d)$ or $H_{CR}(X_W) \otimes \Lambda$. Let $\{T_i\}_{i=0}^n$ be a homogeneous basis of $H$. We choose $T_0$ to be the identity class, i.e. $T_0 = 1_0 \in H_{CR}(X_W)$ in GW theory and $T_0 = \phi_0 \in H(W, \mu_d)_1$ in FJRW theory. We set $g_{ij} = (T_i, T_j)$ and let $(g^{ij})$ denote the matrix inverse to $(g_{ij})$. We write

\[ \left\{ \tau_{b_1}(T_1), \ldots, \tau_{b_n}(T_{n}) \right\}_{0,n}^{FJRW} = \left\{ \tau_{b_1}(T_1), \ldots, \tau_{b_n}(T_{n}) \right\}_{0,n}^{GW} \]  

for FJRW theory; \[ \sum_{\beta \in \mathrm{Eff}} \left\{ \tau_{b_1}(T_1), \ldots, \tau_{b_n}(T_{n}) \right\}_{0,n,\beta}^{GW} Q^\beta \]  

for GW theory.

Let $t^0, \ldots, t^s$ denote the co-ordinates of $H$ dual to the basis $T_0, \ldots, T_s$ such that $t = \sum_{i=0}^s t^i T_i$ denotes a general point\(^{11}\) on $H$. We regard $H$ as a supermanifold such that $t^i$ has the parity $|i| = \deg T_i \mod 2$ and odd co-ordinates anticommute $t^i t^j = (-1)^{|i||j|} t^j t^i$.

Let $K[[t]]$ denote the tensor product of the formal power series ring in even variables and the exterior algebra in odd variables, i.e. $K[[t]] = K[[t^\text{even}]] \otimes_K \bigwedge^* \left( \left( \bigoplus_{|i| \text{odd}} K[t^i] \right) \right)$. The quantum product $\bullet$ is a $K[[t]]$-bilinear product on $H \otimes K[[t]]$ defined by

\[ T_i \bullet T_j = \sum_{k, l=0}^s \sum_{n \geq 0} \frac{1}{n!} (T_i, T_j, T_k, t, \ldots, t)_{0,n+3g^l r} T_l. \]

This is super-commutative and associative by the WDVV equation. We call $(H \otimes K[[t]], \bullet)$ the quantum cohomology ring. The identity of the product $\bullet$ is given by $T_0$. When we set $Q = 0$ and $t = 0$ in Gromov–Witten theory, the product $\bullet$ defines the so-called Chen–Ruan orbifold cup product on $H_{CR}(X_W)$. This limit $Q = t = 0$ is called the large radius limit. On the other hand, the limit point $t = 0$ in FJRW theory is called the Landau–Ginzburg point.

The quantum connection is the set of operators $\nabla_i$, $i = 0, \ldots, s$ on $H \otimes K[[t]][z, z^{-1}]$ defined by

\[ \nabla_i \alpha = \frac{\partial \alpha}{\partial t^i} + \frac{1}{z} T_i \bullet \alpha, \quad \alpha \in H \otimes K[[t]][z, z^{-1}]. \]

Here $z$ is a formal parameter. The associativity of the product $\bullet$ implies that these operators supercommute, i.e. $\nabla_i \nabla_j - (-1)^{|i||j|} \nabla_j \nabla_i = 0$. We regard the quantum cohomology

\[ 10 \text{ Note that the FJRW correlators are not defined for } n = 0, 1, 2 \text{ (since the moduli spaces are empty), but the GW correlators still exist for these } n \text{ because the degree } \beta \text{ can be non-zero. We set the correlator to be zero when the subscript } (0, n) \text{ or } (0, n, \beta) \text{ is not in the stable range.} \]

\[ 11 \text{ We use upper indices for the co-ordinates } t^i. \text{ Note that } t^i \text{ does not mean the } i\text{-th power of } t. \]
$H \otimes K[[t]]$ as the trivial vector bundle over the formal neighbourhood of the origin in $H$ and $\nabla$ as a flat connection on it with parameter $z$. Moreover, we can extend the connection in the $z$-direction because of the homogeneity in FJRW/GW theory. Define a section $E \in H \otimes K[[t]]$ by

$$E := \sum_{i=0}^{s} \left(1 - \frac{1}{2} \deg T_i\right) t^i T_i.$$  

This corresponds to the Euler vector field$^{12}$ $\sum_{i=0}^{s} \left(1 - \frac{1}{2} \deg T_i\right) t^i \frac{\partial}{\partial t_i}$. By abuse of notation, we also denote the vector field by $E$. It defines the half of the grading of variables: $\deg t^i := 2 - \deg T_i$. Let $\text{Gr}$ denote the grading operator

$$\text{Gr}(T_i) := \frac{\deg T_i}{2} T_i.$$  

The connection $\nabla_z$ in the $z$-direction is defined to be

$$\nabla_z \alpha = \frac{\partial \alpha}{\partial z} - \frac{1}{z^2} E \bullet \alpha + \frac{1}{z} \text{Gr}(\alpha).$$  

for $\alpha \in H \otimes K[[t]][z, z^{-1}]$. We have $[\nabla_z, \nabla_i] = 0$, $i = 1, \ldots, s$ and the pairing $(\cdot, \cdot)$ is $\nabla$-flat. The precise meaning of the flatness of $(\cdot, \cdot)$ is as follows. Let $(-)^* : H \otimes K[[t]][z, z^{-1}] \to H \otimes K[[t]][z, z^{-1}]$ denote the map $\alpha(z) \mapsto \alpha(-z)$ flipping the sign of $z$. By extending the pairing $(\cdot, \cdot)$ on $H$ to $H \otimes K[[t]][z, z^{-1}]$ bilinearly over $K[[t]][z, z^{-1}]$, we have

$$\frac{\partial}{\partial t^i} ((-)^* \alpha_1, \alpha_2) = ((-)^* \nabla_i \alpha_1, \alpha_2) + ((-)^* \alpha_1, \nabla_i \alpha_2)$$  

for $\alpha_1, \alpha_2 \in H \otimes K[[t]][z, z^{-1}]$. The flatness of the pairing follows from the Frobenius property $(\alpha_1 \bullet \alpha_2, \alpha_3) = (\alpha_1, \alpha_2 \bullet \alpha_3)$.

We have a canonical solution of the quantum connection. Define $L \in \text{End}(H) \otimes K[[t]][z, z^{-1}])$ by

$$L(t, z) \alpha = \alpha + \sum_{i,j=1}^{s} \sum_{n \geq 0} \sum_{b \geq 0} \frac{1}{n!} (-z)^{b+1}(\tau_i(\alpha), t, \ldots, t, T_i^b)_{i,0,n+b} z^b T_i.$$  

This is an invertible endomorphism satisfying the following differential equations:

**Proposition 2.7.** — For $\alpha \in H$, we have

$$\nabla_t(L(t, z)z^{-\text{Gr}} \alpha) = 0, \quad \nabla_z(L(t, z)z^{-\text{Gr}} \alpha) = 0.$$  

For $\alpha_1, \alpha_2 \in H$, we have

$$(L(t, -z) \alpha_1, L(t, z) \alpha_2) = (\alpha_1, \alpha_2).$$

$^{12}$ Since we are working in the Calabi-Yau case, the term $\mathcal{O}(X)$ vanishes.
Proof. — These are well-known in GW theory (see e.g. [42, Proposition 2.4]) and can be proved similarly for FJRW theory. So we only sketch the outline of the proof in the case of FJRW theory. The equation $\nabla_r (L(t,z)z^{-Gr}t) = 0$ is a formal consequence of the Topological Recursion Relation (TRR), as shown in [54, Proposition 2] for the GW theory. The TRR in FJRW theory is proved in [26, Theorem 4.2.8]. The equation $\nabla_z (L(t,z)z^{-Gr}t) = 0$ follows from the homogeneity (13) of FJRW invariants. Since $L(t,z)\alpha_i$ is flat in the $t$-direction, the pairing $(L(t,-z)\alpha_1, L(t,z)\alpha_2)$ does not depend on $t$ (see (19)). Therefore $(L(t,-z)\alpha_1, L(t,z)\alpha_2) = (L(0,-z)\alpha_1, L(0,z)\alpha_2) = (\alpha_1, \alpha_2)$. □

2.3.1. Restriction to the narrow/ambient part. — Recall that the state space $H$ is decomposed as $H_{nar}(W, \mu_d) \oplus H_{bro}(W, \mu_d)$ for FJRW theory and as $H_{amb}(X_W) \otimes \Lambda \oplus H_{pri}(X_W) \otimes \Lambda$ for GW theory. In this section, we denote this decomposition as

$$H = H' \oplus H''$$

where $H'$ denotes the narrow/ambient part and $H''$ denotes the broad/primitive part. The decomposition is orthogonal with respect to the pairing $(\cdot, \cdot)$. Moreover we have the following.

Proposition 2.8. — For $\alpha_1, \ldots, \alpha_n \in H'$ and $\gamma \in H''$, we have

$$(\tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n), \tau_{c}(\gamma))_{0,n+1} = 0.$$ 

In particular, $H'$ is closed under the quantum product $\cdot$ when the parameter $t$ is restricted to lie on $H'$; the quantum connection $\nabla$ and the fundamental solution $L(t,z)$ preserve $H'$ as far as $t \in H'$.

Proof. — Because of the deformation invariance ([25, Theorem 1.2.5], [10, Theorem 3.4.2]) we can assume that $W$ is of Fermat type $W = x_1^{d/w_1} + \cdots + x_N^{d/w_N}$. Then the maximal diagonal group of symmetries preserving $W$ is given as $G_{max} = \{(z_1, \ldots, z_N) | z_i^{d/w_i} = 1\} \cong \mu_{d/w_1} \times \cdots \times \mu_{d/w_N}$. The $G_{max}$-action on $C^N$ naturally lifts to the state space $H$. We claim that $H'$ is the $G_{max}$-invariant part of $H$:

$$H' = H^{G_{max}}$$

and that $H''$ is the sum of non-trivial irreducible $G_{max}$-representations. The proposition follows from this claim and the $G_{max}$-invariance of the correlator:

$$(\tau_{b_1}(g\alpha_1), \ldots, \tau_{b_n}(g\alpha_n), \tau_{c}(g\gamma))_{0,n+1} = (\tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n), \tau_{c}(\gamma))_{0,n+1}, \quad g \in G_{max}.$$ 

First we show the claim for the FJRW state space. We use the description of $H(W, \mu_d)_k$ as the Jacobi space $\Omega(W_k)^{\mu_d}$ given in (7). If $k \in \text{Nar}$, the sector $H(W, \mu_d)_k$ is obviously $G_{max}$-invariant. Assume that $k \notin \text{Nar}$. Then each element of $H(W, \mu_d)_k$ can
be represented by a sum of monomial \( N_k \)-forms of the form \( \prod_{k_j \in \mathbb{Z}} (x_i^{a_i} d x_i) \) with \( 0 \leq a_i \leq d/w_i - 2 \). But each summand spans a non-trivial irreducible \( G_{\max} \)-representations and the claim follows.

Next we show the claim for the GW state space. Each twisted sector has a decomposition \( \text{H}(\text{P}(w)_f \cap X_W) = \text{H}_{\text{namb}}(\text{P}(w)_f \cap X_W) \oplus \text{H}_{\text{pri}}(\text{P}(w)_f \cap X_W) \). It is obvious that \( \text{H}_{\text{namb}}(\text{P}(w)_f \cap X_W) \) is \( G_{\max} \)-invariant. The primitive part \( \text{H}_{\text{pri}}(\text{P}(w)_f \cap X_W) \) is isomorphic to \( \text{Gr}^W_k \text{H}^{\text{namb}}_{-1}(\text{W}^{-1}(1)) \) for \( k = fd \) (see \([64, \text{p. 216}]\)), which is \( \text{H}(\text{W}, \mu_\sigma)_k \) (see Appendix A). Now the claim follows for the same reason as the previous paragraph, since \( k = fd \notin \text{Nar} \).

**Remark 2.9.** — The fact that the ambient part is closed under the quantum product (in GW theory) is shown \([43, \text{Corollary 2.5}]\) in general for complete intersections.

**2.3.2. Divisor equation and the specialization \( Q = 1 \).** — In orbifold GW theory, we have the following Divisor Equation \([1, \text{Theorem 8.3.1}]\):

\[
\begin{align*}
\langle \tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n), \beta \rangle_{0,n+1,\beta}^{\text{GW}} &= \langle \beta \rangle \langle \tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n) \rangle_{0,n,\beta}^{\text{GW}} \\
&+ \sum_{i,b_i > 0} \langle \tau_{b_i}(\alpha_1), \ldots, \tau_{b_i}(\alpha_n) \rangle_{0,n,\beta}^{\text{GW}},
\end{align*}
\]

where \( \beta = c_1(\mathcal{O}(1)) \) is the hyperplane class on \( X_W \) in the untwisted sector. We choose the homogeneous basis \( \{ T_i \}_{i=0} \) such that \( T_0 = 1_0 \) and also that \( T_1 = \beta \). The divisor equation implies that

\[
T_i \cdot T_j = \sum_{k,l=0}^{s} \sum_{n \geq 0} \sum_{\beta \in \text{Eff}} \langle T_i, T_j, T_k, \ldots, T_{l'} \rangle_{0,n+3,\beta}^{\text{GW}} e^{(\beta,\beta)l} Q^l g^{kl} T_l,
\]

where \( l' = \sum_{i \neq 1} t' T_i \). Therefore, the specialization \( T_i \cdot T_j |_{Q=1} \) makes sense as an element of \( \text{H}_{\text{CR}}(X_W) \otimes \mathbb{C}[[e^{1/w}, t']] \). Here \( w \) is the least common multiple of \( w_1, \ldots, w_N \) (so that \( \mathcal{O}_{\text{P}(w)}(w) \) is a pull-back from the coarse moduli space \( \text{P}(w) \)). Under this specialization, the limit \( e^{1/w} = t' = 0 \) plays the role of the large radius limit. Similarly, the specialization \( Q = 1 \) of the fundamental solution \( L(t, z) \) gives

\[
L(t, z) |_{Q=1} = e^{-1/pz} = \sum_{i,j=1}^{s} \sum_{n \geq 0} \sum_{\beta \in \text{Eff}} \frac{1}{n!(\beta)_t l_1} \times \langle \tau_{b_j}(e^{-1/pz} \alpha), \ldots, t', T_{l'} \rangle_{0,n+2,\beta}^{\text{GW}} e^{(p,\beta)l} g^{p^{l_1} T_j}.
\]

This is an element of \( \text{End}(\text{H}_{\text{CR}}(X_W)) \otimes \mathbb{C}[[e^{1/w}, t']][[z^{-1}]] \) and gives a fundamental solution to the quantum connection \( \nabla |_{Q=1} \). In this way the divisor equation reduces the ground ring from \( \Lambda \) to \( \mathbb{C} \).
2.4. Quantum D-modules and integral structure. — Here we define the narrow/ambient part of the quantum D-module and introduce a certain integral structure on it. In this section we entirely work over \( \mathbb{C} \). Let \( H \) denote the state space over \( \mathbb{C} \) and \( H' \subset H \) denote the narrow/ambient part:

\[
H' = \begin{cases} 
H_{\text{nar}}(W, \mu_d) & \text{for FJRW theory;} \\
H_{\text{amb}}(X_W) & \text{for GW theory.}
\end{cases}
\]

(21)

Recall that \( H' \) is closed under the quantum product by Proposition 2.8. Let \( \{T_0, T_1, \ldots, T_r\} \) \((r \leq s)\) be a homogeneous basis of \( H' \) such that \( T_0 \) is the identity. It is, for example, a reordering of the basis \( \{\phi_{k-1} \mid k \in \text{Nar}\} \) (FJRW theory) or the basis \( \{p^i 1_f \mid i \geq 0, f \in \mathfrak{f}\} \) (GW theory). For GW theory, we choose \( T_1 \) to be the hyperplane class \( p = c_1(\mathcal{O}(1)) \). Let \( \{t^0, \ldots, t^r\} \) denote the dual co-ordinates and \( t = \sum_{i=0}^r t^i T_i \) denote a general point on \( H' \).

The parity of these co-ordinates are all even. In this section, the parameter \( t \) is restricted to lie on \( H' \) unless otherwise stated. Also in GW theory, we set the Novikov parameter \( Q \) to 1 in the quantum product \( \bullet \), the connection \( \nabla \) and the fundamental solutions \( L(t, z) \) as done in Section 2.3.2.

2.4.1. Convergence assumption. — We assume that the quantum product \( T_i \bullet T_j \), \( 0 \leq i, j \leq r \) are all convergent power series. This means

\[
T_i \bullet T_j \in H' \otimes \mathbb{C}\{t^0, t^1, \ldots, t^r\} \quad \text{for FJRW theory;}
\]

\[
T_i \bullet T_j \in H' \otimes \mathbb{C}\{t^0, t^1/w, t^2, \ldots, t^r\} \quad \text{for GW theory.}
\]

Let \( U \subset H' \) denote the domain of convergence of the product \( \bullet \). For FJRW theory, \( U \) is of the form

\[
\{ ||t^i|| < \epsilon, \quad (\forall i) \}
\]

for some \( \epsilon > 0 \). For GW theory, \( U \) is of the form

\[
\{ \Re(t^i) < -M, \quad ||t^i|| < \epsilon, \quad (\forall i \neq 1) \}
\]

for some \( \epsilon, M > 0 \). In practice, we do not need to assume the full convergence of the product. One can consider the quantum D-module over a submanifold of \( U \) where the product \( \bullet \) is convergent. In our case, we show in Section 5.2 that the quantum product \( \bullet \) is convergent on the image of the mirror map. When \( X_W \) is a manifold, we show the full convergence in Section 5.5 for both GW and FJRW theories.

Note that the convergence assumption imply that \( \nabla \) and \( L(t, z) \) are analytic in \( t \in U \) and \( z \in \mathbb{C}^\times \).
2.4.2. Quantum D-module. — Let $U \subset H'$ be as in Section 2.4.1. The present quantum D-module is defined as a meromorphic flat connection over $U \times \mathbb{C}$. Let $z$ denote the co-ordinate on the second factor $\mathbb{C}$ and $\pi : U \times \mathbb{C} \to U$ denote the projection to the first factor. Let $(-) : U \times \mathbb{C} \to U \times \mathbb{C}$ be the map sending $(t, z)$ to $(t, -z)$.

**Definition 2.10.** — Let $F = H' \times (U \times \mathbb{C}) \to U \times \mathbb{C}$ be the trivial vector bundle with fibre $H'$. Let $\nabla$ be the meromorphic flat connection (quantum connection) on $F$.

$$\nabla = d + \frac{1}{z} \sum_{i=0}^r (T_i \cdot t^i) + \left( -\frac{1}{z} (E \cdot t) + Gr \right) \frac{dz}{z}$$

which can be regarded as a map

$$\nabla : \mathcal{O}(F) \to \mathcal{O}(F)(U \times \{0\}) \otimes \left( \pi^* \Omega^1_U \oplus \mathcal{O}_{U \times \mathbb{C}} \frac{dz}{z} \right).$$

Here $\mathcal{O}(F)(U \times \{0\})$ denotes the sheaf of holomorphic sections of $F$ with at most simple poles along $\{z = 0\} = U \times \{0\}$. Let $P$ be an $\mathcal{O}_{U \times \mathbb{C}}$-bilinear pairing

$$((-)^* s_1, s_2) := (2\pi \hat{c} z)^{\frac{\hat{c}}{2}} (s_1(t, -z), s_2(t, z)).$$

Here $\hat{c} = \dim X_W = N - 2$ and $(\cdot, \cdot)$ in the right-hand side denotes the standard pairing on the state space defined in (4) and (14). The pairing satisfies

$$(-1)^{\hat{c}} P((-)^* s_1, s_2) = (-)^* P((-)^* s_2, s_1)$$

$$dP((-)^* s_1, s_2) = P((-)^* \nabla s_1, s_2) + P((-)^* s_1, \nabla s_2).$$

We call the tuple $(F, \nabla, P)$ the narrow part quantum D-module $QDM_{nar}(W, \mu_d)$ (in the case of $F\bar{J}RW$ theory) and the ambient part quantum D-module $QDM_{amb}(X_W)$ (in the case of GW theory).

**Remark 2.11.** — In [49] and [42, Definition 2.2], the quantum connection $\nabla_{\varphi}$ in the $z$-direction is shifted by $-\hat{c}/2$ so that it makes the ordinary pairing $P/(2\pi \hat{c} z)^{\hat{c}}$ flat. In this paper, we adopt the convention in [43, Definition 3.1] because it is more compatible with mirror symmetry.

**Remark 2.12.** — The quantum D-module here is a (TEP$(\hat{c})$) structure in the sense of Hertling [35, Remark 2.13]. Moreover, by the given trivialization, $F$ is extended over $U \times \mathbb{P}^1$ as a trivial vector bundle and thus gives a (trTLEP$(\hat{c})$)-structure [35, Definition 5.5]. In the context of LG/CY correspondence, it is more convenient to consider the connection over $U \times \mathbb{C}$ since the extensions across $z = \infty$ do not match under analytic continuation.
Remark 2.13. — Over the degree two subspace \( H'^2 \subset H' \), the narrow/ambient part quantum D-module gives rise to a variation of Hodge structure, the so-called (narrow/ambient) A-model VHS [43, Section 6.2]. This is defined to be the restriction of \((\mathcal{F}, \nabla)\) to the subspace \((H'^2 \cap U) \times \{ z = 1 \}\) equipped with the decreasing Hodge filtration

\[
\mathcal{F}^p \colonequals H'^{2(p-2)} \otimes \mathcal{O}_{H'^2 \cap U}
\]

and the polarization

\[
Q_A(\alpha, \beta) = (2\pi i)^p \left( (-1)^{\text{deg} \alpha} \mathcal{O}_{H'^2 \cap U} \right) \mathcal{O}_{H'^2 \cap U}
\]

2.4.3. Galois action. — The quantum D-module has an important discrete symmetry which we call the Galois action. This symmetry is also compatible with mirror symmetry.

Proposition 2.14 (Galois action in FJRW theory). — Let \( H \) be the FJRW state space and \( H' \) be its narrow part. Define the linear map \( G : H' \to H' \) by

\[
G|_{H(W, \mu_k)} = e^{-2\pi i (k-1)/d} \text{id}_{H(W, \mu_k)}.
\]

The map \( G \) preserves \( H' \). Without loss of generality, one can assume that the convergence domain \( U \subset H' \) is preserved by \( G \). The bundle map \( G_F : F \to G^* F \) defined by

\[
G_F : H' \times (U \times \mathbb{C}) \to H' \times (U \times \mathbb{C})
\]

\[
(\alpha, (t, z)) \mapsto (e^{-2\pi i/d} G(\alpha), (G(t), z))
\]

preserves the connection \( \nabla \) (i.e. \( \nabla_{G(v)} \circ G_F = G_F \circ \nabla_v \)) and the pairing \( \mathcal{P} \). We call it the Galois action of the narrow part quantum D-module.

Proof. — For a \( d \)-spin structure \( L \to C \) on a pointed orbicurve \((C, \sigma_1, \ldots, \sigma_n)\), we have \( \text{deg}(L) - \sum_{i=1}^n \text{age}_{\sigma_i}(L) \in \mathbb{Z} \) by Riemann-Roch for orbicurves [1, 45, 65]. Thus the moduli space \( \text{Spin}^d_{0,n}(k_1, \ldots, k_n) \) is empty unless \( 2 + \sum_{i=1}^n k_i \equiv 0 \mod d \). Therefore we have

\[
\left\{ \tau_{k_1} \left( G(\phi_{k_1}) \right), \ldots, \tau_{k_n} \left( G(\phi_{k_n}) \right) \right\}_{FJR} = e^{2(2\pi i)/d} \left\{ \tau_{k_1} \left( \phi_{k_1} \right), \ldots, \tau_{k_n} \left( \phi_{k_n} \right) \right\}_{FJR}
\]

for \( k_1 + 1, \ldots, k_n + 1 \in \text{Nar} \). This fact and the formula (5) of the pairing show that

\[
G(\alpha_1) \bullet_{G(v)} G(\alpha_2) = G(\alpha_1 \bullet \alpha_2)
\]

for \( \alpha_1, \alpha_2 \in H' \). Here the subscripts of \( \bullet \) denote the parameter of the quantum product. The statement follows easily from this. \( \square \)
Remark 2.15. — Since the FJRW invariants in our case are (regardless of narrow or broad) certain intersection numbers on $\text{Spin}_{0,n}^d(k_1, \ldots, k_n)$, the same argument shows that the Galois action preserves $\nabla$ and $P$ defined on the full FJRW state space $H$.

Proposition 2.16 (Galois action in GW theory: [42, Proposition 2.3]). — Let $H$ be the GW state space and $H'$ be its ambient part. Define $G : H \to H$ to be the affine-linear map

$$G(\alpha) = e^{2\pi i \alpha} - 2\pi i \rho, \quad \text{for} \ \alpha \in H(P(w) \cap X_W), \quad f \in \mathfrak{g}.$$  

The map $G$ preserves $H'$. Without loss of generality, one can assume that the convergence domain $U \subset H'$ is preserved by $G$. The bundle map $G_F : F \to G^*F$ defined by

$$H' \times (U \times C) \to H' \times (U \times C)$$

$$(\alpha, (t, z)) \mapsto (\alpha G(\alpha), (G(t), z))$$

preserves the connection $\nabla$ and the pairing $P$. Here $\partial G$ is the differential (linear part) of $G$. We call it the Galois action of the ambient part quantum D-module.

Proof. — In [42, Proposition 2.3], the Galois action was defined for each orbifold line bundle. The map $G_F$ here arises from $\mathcal{O}(1)$. □

Via the Galois action, the quantum D-module $(F, \nabla, P)$ over $U \times C$ descends to a flat connection on the quotient space $(U/(G)) \times C$. We denote it by $(F, \nabla, P)/\langle G \rangle$.

2.4.4. Integral structure. — The $\hat{\Gamma}$-integral structure in the orbifold GW theory was introduced in [42, Section 2.4], [44, Section 3.1]. We generalize it to the case of FJRW theory for $(\mathbb{C}^N, W, \mu_d)$.

Definition 2.17.

1. In FJRW theory, the Gamma class $\hat{\Gamma}_{\text{FJRW}} \in \text{End}(H)$ is defined to be

$$\hat{\Gamma}_{\text{FJRW}} := \bigoplus_{k=0}^{d-1} \prod_{j=1}^N \Gamma(1 - \langle kq_j \rangle)$$

where $q_j = w_j/d$ and the $k$th summand acts on $H(W, \mu_d)_k$ by the scalar multiplication. This is similar to the Gamma class [42, Section 2.4] of the tangent bundle of the orbifold $[\mathbb{C}^N/\mu_d]$.

2. In GW theory, the Gamma class $\hat{\Gamma}_{\text{GW}} \in \text{End}(H)$ is defined to be

$$\hat{\Gamma}_{\text{GW}} := \bigoplus_{f \in \mathfrak{g}} \prod_{i=1}^N \Gamma(1 - \langle f w_i \rangle + w_i \rho) \cdot \Gamma(1 + \rho)$$
where \( p = c_1(O(1)) \) is the hyperplane class and the summand indexed by \( f \in \mathfrak{F} \) acts on \( H(\mathbb{P}(w)_f \cap X_W) \) by the cup product. This is the Gamma class [42, Section 2.4] of the tangent bundle of the Calabi-Yau hypersurface \( X_W \).

**Remark 2.18.** — Libgober [48] observed that the Gamma class arises from periods of mirrors of Calabi-Yau hypersurfaces.

We introduce a flat section associated to a graded matrix factorization of \( W \) (see Section 4.1) or a vector bundle on \( X_W \). We use the Chern character map

\[
\text{ch} : \text{MF}_{\mu_d}(W) \to \bigoplus_{k=0}^{d-1} \Omega(W)^{\mu_d} \cong H(W, \mu_d) \quad \text{for FJRW theory};
\]

\[
\text{ch} : D^b(X_W) \to H_{CR}(X_W) \quad \text{for GW theory}.
\]

The Chern character for a matrix factorization (due to Walcher [68] and Polishchuk-Vaintrob [56]) will be explained in Section 4.1.1 and we use the isomorphism \( \bigoplus_{k=0}^{d-1} \Omega(W)^{\mu_d} \cong H(W, \mu_d) \) in Proposition 2.1. The Chern character for a orbi-vector bundle is the “stabilizer-equivariant” version which appears in the Kawasaki-Riemann-Roch formula [45, 65]. For \( \text{ch} : D^b(X_W) \to H_{CR}(X_W) \), see for instance [42, Section 2.4] where it is denoted by \( \tilde{\text{ch}} \).

**Definition 2.19** (\( \hat{\Gamma} \)-integral structure). — Let \( \deg_0 : H \to H \) be the degree operator without the shift (“bare” degree operator):

\[
\deg_0 := -2 \text{id}_{H(W, \mu_d)} \quad \text{for FJRW theory};
\]

\[
\deg_0 |_{H^*(\mathbb{P}(w)_f \cap X_W)} := n \text{id}_{H^*(\mathbb{P}(w)_f \cap X_W)} \quad \text{for GW theory}.
\]

Let \( \text{inv} : H \to H \) denote the map induced from the natural isomorphisms

\[
H(W, \mu_d)_k \cong H(W, \mu_d)_{d-k} \quad \text{for FJRW theory};
\]

\[
H(\mathbb{P}(w)_f \cap X_W) \cong H(\mathbb{P}(w)_{(1-f)} \cap X_W) \quad \text{for GW theory}.
\]

Let \( E \) be an object of \( \text{MF}_{\mu_d}(W) \) (in the case of FJRW theory) or an object of \( D^b(X_W) \) (in the case of GW theory). We define a \( \hat{\Gamma} \)-flat section \( s(E) \) by

\[
s(E)(t, z) := \frac{1}{(2\pi i)^N} L(t, z) z^{-\hat{\Gamma} \cdot \text{inv} \ast \text{ch}(E)}
\]

where \( \hat{\Gamma}, L(t, z) \) and \( \text{ch}(E) \) are the Gamma class, the fundamental solution (20) and the Chern character in the respective theory. It is clear from the definition that \( s(E) \) depends only on the numerical class of \( E \).
When \( \text{ch}(\mathcal{E}) \in H^* \), \( s(\mathcal{E}) \) defines a flat section of narrow/ambient part quantum D-module. Define the \( \mathbb{Z} \)-local system over \( U \times \mathbb{C}^x \) by

\[
F_{\mathbb{Z}} := \{ s(\mathcal{E}) \mid \text{ch}(\mathcal{E}) \in H' \} \subset \Gamma(U \times \mathbb{C}^x, \mathcal{O}(F))^\nabla
\]

where \( \mathcal{E} \) ranges over objects of \( \text{MF}_{\mu_d}^\text{gr}(W) \) or \( D^b(X_W) \). (Note that \( \text{ch}(\mathcal{E}) \) lies in \( H \), but not in \( H' \) in general.) We call this the \( \hat{\Gamma} \)-integral structure of the narrow/ambient part quantum D-module.

**Remark 2.20.** — The degree \( \deg_0 \) is even on the image of the Chern character map. In fact, the Chern character takes values in the \( \hat{\Gamma} \)-local system over \( U \). Later, in Part (1), since the Galois action preserves \( \nabla \), it suffices to check the equality at \( t = 0 \) (see Remark 3.6). For Part (1), since the Galois action preserves \( \nabla \), it suffices to check the equality at \( t = 0 \) (see Remark 3.6). For Part (1), since the Galois action preserves \( \nabla \), it suffices to check the equality at \( t = 0 \) (see Remark 3.6). For Part (1), since the Galois action preserves \( \nabla \), it suffices to check the equality at \( t = 0 \) (see Remark 3.6).

**Proposition 2.21.**

1. The \( \hat{\Gamma} \)-integral structure is preserved under the Galois action, i.e.

\[
e^{-2\pi i/d} G(s(\mathcal{E}))(G^{-1}(t), z) = s(\mathcal{E}(1))(t, z) \quad \text{for FJRW theory},
\]

\[
e^d G(s(\mathcal{E}))(G^{-1}(t), z) = s(\mathcal{E} \otimes \mathcal{O}(-1))(t, z) \quad \text{for GW theory},
\]

where \( \mathcal{E}(1) \) is the shift of the grading of \( \mathcal{E} \) by 1. In particular, \( F_{\mathbb{Z}} \) defines an integral structure on the quotient \( (F, \nabla, \mathcal{P})/\langle G \rangle \).

2. We have

\[
P((-)^*s(\mathcal{E}), s(\mathcal{F})) = (-1)^{N-1} \chi(\mathcal{E}, \mathcal{F}) \quad \text{for FJRW theory},
\]

\[
P((-)^*s(\mathcal{E}), s(\mathcal{F})) = \chi(\mathcal{E}, \mathcal{F}) \quad \text{for GW theory},
\]

where \( \chi(\mathcal{E}, \mathcal{F}) := \sum_{i \in \mathbb{Z}} (-1)^i \dim \text{Hom}(\mathcal{E}, \mathcal{F}[i]) \) is the Euler pairing of \( \text{MF}_{\mu_d}^\text{gr}(W) \) or \( D^b(X_W) \). In particular, \( P \) takes values in \( \mathbb{Z} \) on \( F_{\mathbb{Z}} \).

**Proof.** — The proof relies on the Hirzebruch-Riemann-Roch formula in each category. We explain the case of FJRW theory. See [42, Proposition 2.10] (or [43, Definition 3.6]) for the case of GW theory. For Part (1), since the Galois action preserves \( \nabla \), it suffices to check the equality at \( t = 0 \) (see Remark 2.15). It follows from \( L(0, z) = \text{id} \) and

\[
e^{-2\pi i/d} G(\text{inv}^* \text{ch}(\mathcal{E})) = \text{inv}^* \text{ch}(\mathcal{E}(1)).
\]

Next we show Part (2). Setting \( \Psi(\mathcal{E}) = \hat{\Gamma}((2\pi i)^{\deg} \text{inv}^* \text{ch}(\mathcal{E})) \), we have

\[
P((-)^*s(\mathcal{E}), s(\mathcal{F}))
\]

\[
= (2\pi i)^{-i} \hat{\Gamma}(((-z)^{-\text{Gr}} \Psi(\mathcal{E}), x^{-\text{Gr}} \Psi(\mathcal{F})) \quad \text{by Proposition 2.7}
\]

\[
= (2\pi i)^{-i} \hat{\Gamma}((-1)^{-\text{Gr}} \Psi(\mathcal{E}), \Psi(\mathcal{F}))
\]

\[
= (2\pi i)^{-i} \sum_{k=0}^{d-1} \left( \prod_{j=1}^{N} \Gamma(1 - (kj)) \Gamma(1 - (d - k,j)) \right) (2\pi i)^{-2}
\]

Next we show Part (2). Setting \( \Psi(\mathcal{E}) = \hat{\Gamma}((2\pi i)^{\deg} \text{inv}^* \text{ch}(\mathcal{E})) \), we have

\[
P((-)^*s(\mathcal{E}), s(\mathcal{F}))
\]

\[
= (2\pi i)^{-i} \hat{\Gamma}(((-z)^{-\text{Gr}} \Psi(\mathcal{E}), x^{-\text{Gr}} \Psi(\mathcal{F})) \quad \text{by Proposition 2.7}
\]

\[
= (2\pi i)^{-i} \hat{\Gamma}((-1)^{-\text{Gr}} \Psi(\mathcal{E}), \Psi(\mathcal{F}))
\]

\[
= (2\pi i)^{-i} \sum_{k=0}^{d-1} \left( \prod_{j=1}^{N} \Gamma(1 - (kj)) \Gamma(1 - (d - k,j)) \right) (2\pi i)^{-2}
\]
where \( \text{ch}(\mathcal{F})_k \) denotes the component of \( \text{ch}(\mathcal{F}) \) in the sector \( H(W, \mu_d)_k \) and we used the equality \( \Gamma(1 - x)\Gamma(x) = \pi / \sin(\pi x) \). Note that \( \text{ch}(\mathcal{F})_k \) vanishes if \( N_k \) is odd. By Proposition 2.1, we can write the last expression in terms of the residue pairing:

\[
\sum_{k=0}^{d-1} \left( \prod_{\langle k \rangle \neq 0} \frac{1}{1 - e^{2\pi i \langle k \rangle}} \right) (-1)^{N-1} (-1)^{-\frac{N(N_l-1)}{2}} \frac{1}{d} \text{Res}_{W_k} (\text{ch}(\mathcal{E})_{d-k}, \text{ch}(\mathcal{F})_k)
\]

where we used the fact that the degree of \( \text{ch}(\mathcal{F})_k \) as an element of \( \Omega(W_k)^{\mu_d} \) is \( (N_k/2)d \) (see Remark 4.4; the degree here is different from the degree as an element of the FJRW state space). This equals \( (-1)^{N-1} \chi(\mathcal{E}, \mathcal{F}) \) by Hirzebruch-Riemann-Roch Theorem 4.6.

**Remark 2.22.** — In Proposition 2.21, we do not need to assume that \( \text{ch}(\mathcal{E}) \in H' \) or \( t \in H' \).

2.5. **Statements of the main results.** — In this section we summarize our main results in three theorems. Theorems 2.23, 2.25 are about analytic continuation of quantum D-modules (with integral structure) and Theorem 2.26 is about the monodromy representation and derived equivalences. These theorems are more precise versions of Theorems 1.1 and 1.2 in the introduction. The proofs take the entire paper and are completed in Section 5 (see Section 2.6).

Let \( v \) be an inhomogeneous co-ordinate of \( \mathbf{P}(1, d) \) such that \( v = \infty \) is the \( \mu_d \)-orbifold point. Then \( u = v^{-1/d} \) gives a uniformizing co-ordinate around the orbifold point (LG point). Set \( \mathcal{M} := \mathbf{P}(1, d) \setminus \{ v = 0, v = v_c \} \), where \( v = 0 \) is the large radius limit point and \( v = v_c := d^{-d} \prod_{j=1}^{N} w_j \) is the conifold point. We write \( (-) : \mathcal{M} \times \mathbf{C}_z \rightarrow \mathcal{M} \times \mathbf{C}_z \) for the map sending \((v, z)\) to \((v, -z)\). The following theorem is a precise version of Theorem 1.1 (without the part concerning the Orlov equivalence). The proof will be completed in Section 5.4.

**Theorem 2.23.** — There exists a locally free sheaf \( \mathcal{F} \) over \( \mathcal{M} \times \mathbf{C}_z \) with a meromorphic flat connection \( \nabla \) (with simple poles along \( z = 0 \))

\[
\nabla : \mathcal{F} \rightarrow \mathcal{F}(\mathcal{M} \times \{0\}) \otimes \Omega_{\mathcal{M} \times \mathbf{C}_z}^1,
\]

a \( \nabla \)-flat, symmetric and non-degenerate pairing

\[
P : (-)^* \mathcal{F} \otimes \mathcal{F} \rightarrow \xi^* \mathcal{O}_{\mathcal{M} \times \mathbf{C}_z}
\]
and a \(\mathbb{Z}\)-local subsystem \(F_z\) of the same rank over \(\mathcal{M} \times \mathbb{C}^*_z\)

\[
F_z \subset (\mathcal{F}|_{\mathcal{M} \times \mathbb{C}^*_z})^V
\]

such that the following holds:

(i) For a small open neighbourhood \(U_{\text{FJRW}} = \{|u| < \epsilon\} \subset \mathcal{M}\) of the LG point, we have a mirror map \(\tau_{\text{FJRW}} : U_{\text{FJRW}} \to H^2_{\text{nar}}(W, \mu_d)/\langle G \rangle\) such that \(\tau_{\text{FJRW}} = -u\phi_1 + O(u^2)\) and

\[
(F, \nabla, (-1)^{N-1}P, F_z)\big|_{U_{\text{FJRW}} \times \mathbb{C}} \cong \tau_{\text{FJRW}}^*(\text{QDM}_{\text{nar}}(W, \mu_d)/\langle G \rangle),
\]

where in the right-hand side appears the narrow part quantum D-module (Definition 2.10) of \((\mathbb{C}^N, W, \mu_d)\) equipped with the \(\hat{\Gamma}\)-integral structure (Definition 2.19) and \(G\) is the Galois action (Section 2.4.3).

(ii) For a small open neighbourhood \(U_{\text{GW}} = \{|v| < \epsilon\} \subset \mathcal{M}\) of the large radius limit point, we have a mirror map \(\tau_{\text{GW}} : U_{\text{GW}} \to H^2_{\text{amb}}(X_W)/\langle G \rangle\) such that \(\tau_{\text{GW}}(v) = p \log v + O(v)\) and

\[
(F, \nabla, P, F_z)\big|_{U_{\text{GW}} \times \mathbb{C}} \cong \tau_{\text{GW}}^*(\text{QDM}_{\text{amb}}(X_W)/\langle G \rangle),
\]

where in the right-hand side appears the ambient part quantum D-module of \(X_W\) equipped with the \(\hat{\Gamma}\)-integral structure and \(G\) is the Galois action.

Remark 2.24. — Restricting the global D-module \((F, \nabla, P, F_z)\) to \(z = 1\), we obtain the analytic continuation between the narrow A-model VHS of \((\mathbb{C}^N, W, \mu_d)\) and the ambient A-model VHS of \(X_W\) in Remark 2.13. The fibre \(F_{(x,1)}\) at \((x, 1) \in \mathcal{M} \times \mathbb{C}_z\) has a well-defined filtration and a polarization

\[
F^\#(F_{(x,1)}) = \{v \in F_{(x,1)} \mid s_v(z) \text{ has a pole of order } \leq \hat{c} - \rho \text{ at } z = 0\}
\]

\[
Q(v_1, v_2) = P(s_{v_1}(-1), s_{v_2}(1)), \quad v_1, v_2 \in F_{(x,1)}
\]

where \(s_v(z) \in H^0(\mathbb{C}^*_z, F|_{(t) \times \mathbb{C}^*_z})\) is a unique \(\nabla\)-flat section such that \(s_v(1) = v\). The filtration and the polarization coincide with those of the A-model VHS near the respective limit point. By analytic continuation, we have an isomorphism of state spaces

\[
\Theta(x) : (H_{\text{nar}}(W, \mu_d), F^\#_\Lambda, Q_\Lambda) \cong (H_{\text{amb}}(X_W), F^\#_\Lambda, Q_\Lambda)
\]

for a point \(x\) on the universal cover \(\tilde{\mathcal{M}}\). Taking the associated graded vector space with respect to \(F^\#\), we can turn this into a graded isomorphism (preserving the polarization). Note that \(\Theta\) does not map the identity to the identity because of the factor \(F\) in the asymptotics (74).
In the case where the Calabi-Yau hypersurface $X_W$ is a smooth manifold (e.g. $P(w) = P^n$ or $P(1, 1, 1, 2)$, $P(1, 1, 1, 1, 1, 2, 3)$, etc.), we can use the reconstruction theorem to prove that the “big” quantum D-modules are analytically continued to each other. Here the word “big” means the quantum D-module over the full narrow/ambient sector $H'$. This is used in contrast with the “small” quantum D-module which is the restriction of the big one to the image of the mirror maps. The following theorem will be proved in Section 5.5.

**Theorem 2.25.** — Assume that $X_W$ is a manifold.

(i) The “big” quantum product of $(C^N, W, \mu, d)$ on the narrow part and the “big” quantum product of $X_W$ on the ambient part are convergent in the sense of Section 2.4.1.

(ii) The global D-module $(F, \nabla, P, F_Z)$ in Theorem 2.23 can be extended to a D-module $(F_{\text{ext}}, \nabla_{\text{ext}}, P_{\text{ext}}, F_Z)$ over a base $M_{\text{ext}} \times C_z$, where $M_{\text{ext}}$ is a complex manifold of dimension $\text{rank} F = \dim H_{\text{amb}}(X_W)$ which contains a Zariski open subset $M'$ of $M$ as a submanifold. The extended D-module is identified with the “big” narrow/ambient part quantum D-module of FJRW/GW theory in a neighbourhood of $U_{\text{GW}}$ or $U_{\text{FJRW}}$.

More precisely, there exists a locally free sheaf $F_{\text{ext}}$ over $M_{\text{ext}} \times C_z$ equipped with a meromorphic flat connection $\nabla_{\text{ext}}$ (with poles of order two along $z = 0$)

$\nabla_{\text{ext}} : F_{\text{ext}} \to F_{\text{ext}}(\mathcal{M}_{\text{ext}} \times \{0\}) \otimes \left(\pi^* \Omega^1_{\mathcal{M}_{\text{ext}}} \oplus \mathcal{O}_{\mathcal{M}_{\text{ext}} \times C_z} \frac{dz}{z}\right)$,

where $\pi : \mathcal{M}_{\text{ext}} \times C_z \to \mathcal{M}_{\text{ext}}$ is the projection, a $\nabla_{\text{ext}}$-flat, symmetric and non-degenerate pairing

$P_{\text{ext}} : (-)^* F \otimes F \to \hat{\mathcal{O}}_{\mathcal{M}_{\text{ext}} \times C_z}$

and a $\mathbb{Z}$-local subsystem $F_Z^{\text{ext}}$ of the same rank over $M_{\text{ext}} \times C_z$

$F_Z^{\text{ext}} \subset (F_{\text{ext}}|_{M_{\text{ext}} \times C_z})^{\nabla_{\text{ext}}}$

such that the following holds:

- $(F_{\text{ext}}, \nabla_{\text{ext}}, P_{\text{ext}}, F_Z^{\text{ext}})|_{\mathcal{M}} = (F, \nabla, P, F_Z)|_{\mathcal{M}}$;
- There exist open neighbourhoods $U_{\text{GW}}$ of $U$ in $\mathcal{M}_{\text{ext}}$ and open embeddings

$\tau_{\text{GW}} : U_{\text{GW}} \hookrightarrow H/\langle G \rangle, \quad \tau_{\text{FJRW}}|_{U_{\text{FJRW}}} = \tau_{\text{GW}}$

for $\bowtie = \text{GW}$ and $\text{FJRW}$ such that we have isomorphisms

$(F_{\text{ext}}^{\text{ext}}, \nabla_{\text{ext}}, (-1)^{N-1}P_{\text{ext}}, F_Z^{\text{ext}})|_{U_{\text{FJRW}}}$

$\cong \tau_{\text{FJRW}}^* (\text{QDM}_{\text{nar}}(W, \mu, d)/\langle G \rangle)$.
Finally we give a statement on the monodromy representation of the global quantum D-module $(\mathcal{F}, \nabla, P, F_Z)$. We choose base points $b_0, b_\infty \in \mathcal{M}$ near the large radius limit point and the LG point such that $b_0, b_\infty \in \mathbb{R}_{>0}$ and $0 < b_0 \ll 1 \ll b_\infty$. We choose paths $\gamma_{\text{CY}}, \gamma_{\text{LG}}, \gamma_0, \gamma_1$ in $\mathcal{M}$ as in Figure 2. We also define (see also Figure 3)

$$
\gamma_l := \gamma_{\text{LG}} \circ \gamma_0 \circ \gamma_{\text{CY}}', \quad \gamma_{\text{con}} := \gamma_{-1}^{-1} \circ \gamma_0
$$

for $l \in \mathbb{Z}$. We adopt the convention that the composite $\gamma_A \circ \gamma_B$ means the concatenation of $\gamma_A$ at the end of $\gamma_B$.

Let $N(X_W)$ denote the numerical $K$-group of $X_W$. From the definition of the $\tilde{\Gamma}$-integral structure, the fibre at $b_0$ of the global $\mathbb{Z}$-local system $F_Z$ is identified with

$$
N'(X_W) := \{ E \in N(X_W) \mid \text{ch}(E) \in H_{\text{amb}}(X_W) \}.
$$

Similarly, the fibre at $b_\infty$ of $F_Z$ is identified with the group $N'(W, \mu_d)$ of numerical classes of matrix factorizations $\mathcal{E}$ such that $\text{ch}(\mathcal{E}) \in H_{\text{nat}}(W, \mu_d)$. The following theorem is a detailed version of Theorem 1.2 (which also includes the part of Theorem 1.1 concerning the Orlov equivalence). The proof will be given in Section 5.6.

**Theorem 2.26.** — The $\mathbb{Z}$-local system $F_Z$ of the global D-module $(\mathcal{F}, \nabla, P, F_Z)$ in Theorem 2.23 induces the representation of the quiver of Figure 2 given by the assignment $b_0 \mapsto N'(X_W)$, $b_\infty \mapsto N'(W, \mu_d)$ and

$$
\begin{align*}
\gamma_{\text{CY}} & \mapsto \mathcal{O}(-1) : N'(X_W) \to N'(X_W) \\
\gamma_{\text{LG}} & \mapsto (1) : N'(W, \mu_d) \to N'(W, \mu_d) \\
\gamma_{-1}^{-1} & \mapsto \Phi_l : N'(W, \mu_d) \to N'(X_W) \\
\gamma_{\text{con}}^{-1} & \mapsto T_{\mathcal{O}} : N'(X_W) \to N'(X_W)
\end{align*}
$$

where $\mathcal{O}(-1)$ denotes the tensor product by $\mathcal{O}(-1)$, $(1)$ denotes the shift of the grading by 1, $\Phi_l$ denotes the Orlov equivalence (Section 4.2) defined for $l \in \mathbb{Z}$ and $T_{\mathcal{O}}$ denotes the Seidel-Thomas spherical twist (Section 5.6) by the structure sheaf. Moreover, the monodromy representation

$$
\rho : \pi_1(M, b_0) \to \text{Aut}(N'(X_W), \chi)
$$
can be lifted to a group homomorphism
\[ \hat{\rho} : \pi_1(\mathcal{M}, b_0) \to \text{Auteq}(\mathcal{D}^b(X_W))/[2], \]
where \([2]\) is the 2-shift functor.

2.6. Flowchart of the proof: — Here we show by a picture the organization of the rest of the paper. Several key results are highlighted in the chart.

The essence of the paper lies in the computation in Section 4 where the analytic continuation map \(U\) of the two I-functions (or more precisely the two \(\mathcal{H}\)-functions) is matched up with the Orlov derived equivalence.

3. Computing FJRW theory

We compute FJRW invariants attached to narrow state space entries. In Section 3.1, we provide an extension of the definition of the invariants to a larger state space. The new invariants are zero on the extended part, but arise as the non-equivariant limit of the \(e_T\)-twisted invariants. In Sections 3.2–3.4, we calculate the twisted invariants (or more precisely the I-function) using Chiodo-Zvonkine’s results [16] and Givental’s symplectic formalism [29].
3.1. Extending FjRW theory. — Define the extended narrow state space (or simply the extended state space) to be

\[ H_{\text{ext}} = \bigoplus_{k=1}^{d} \phi_{k-1}C = H_{\text{nar}}(W, \mu_d) \oplus \bigoplus_{k \notin \text{Nar}} \phi_{k-1}C. \]

This modified state space may be regarded as the result of the replacement of each term of the broad sector \( H(W, \mu_d)_k, k \notin \text{Nar} \), with a one-dimensional term \( \phi_{k-1}C \). As we show straight away in Proposition 3.2 these new states play the role of placeholders in the theory: they only yield vanishing invariants and they allow to simplify the computation of the invariants with narrow entries.

We need to extend the grading of \( H_{\text{nar}}(W, \mu_d) \) to the extended state space. We set (cf. (1))

\[ \deg \phi_{k-1} = 2 \sum_{j=1}^{N} ((k-1)q_j) = 2N_k + 2 \sum_{j=1}^{N} \langle kq_j \rangle - 2 \]

with \( q_j := w_j/d \). The relevant moduli stack is \( \text{Spin}^d_{0,a}(k_1, \ldots, k_n) \) defined as in (9), but for \( k_1, \ldots, k_n \in \{0, \ldots, d-1\} \). Its universal curve \( \pi : C \to \text{Spin}^d_{0,a}(k_1, \ldots, k_n) \) is equipped with a universal \( d \)-spin structure \( \mathcal{L} \) and a line bundle \( \mathcal{M}_i = \mathcal{O}(D_i) \), where \( D_i \subset C \) denotes the divisor of the \( i \)th marking. The extended obstruction \( K \)-class is defined to be

\[ -R\pi_* \left( \bigoplus_{j=1}^{N} \mathcal{L}^{w_j} \right) \quad \text{for} \quad \mathcal{L} = \mathcal{L} \otimes \mathcal{M}^\vee \]

where \( \mathcal{M} = \bigotimes_{i=1}^{n} \mathcal{M}_i \). Let \( p : C \to \mathcal{C} \) denote the morphism forgetting the stack-theoretic structure along all the markings \( D_1, \ldots, D_n \) (but not along the nodes). Then we have

\[ \text{age}_{D_i}(\mathcal{L}) = \frac{k_i}{d}, \quad \mathcal{L}^{\otimes d} \cong p^*\mathcal{w} \]

for the relative dualizing sheaf \( \mathcal{w} \) of \( \pi : \mathcal{C} \to \text{Spin}^d_{0,a}(k_1, \ldots, k_n) \).

Proposition 3.1. — For any fibre \( C \) of \( \mathcal{C} \), we have \( H^0(C, \mathcal{L}^{\otimes w_j}|_C) = 0, j = 1, \ldots, N \). As a consequence, \( R^1\pi_*(\mathcal{L}^{\otimes w_j}) \) is locally free and the extended obstruction \( K \)-class is represented by a vector bundle.

Proof. — Because \( w_j \) divides \( d \), \( \mathcal{L}^{\otimes w_j} \) is a root of \( p^*\mathcal{w} \). On the other hand, we have \( H^0(C, p^*\mathcal{w}) = H^0(\mathcal{C}, \mathcal{w}) = 0 \) because the genus of \( \mathcal{C} = p(C) \) is zero. Hence \( \mathcal{L}^{\otimes w_j}|_C \) does not have nonzero global sections either. \( \Box \)
We define the extended FJRW invariants to be

\[
\tau_{b_i}(\phi_{b_i}), \ldots, \tau_{b_a}(\phi_{b_a})_{\text{FJRW,ext}} := \int_{[\text{Spin}^c_{b_i}(\mathbb{C})]} \left( \prod_{i=1}^n \psi_{i}^{d_i} \right) \cup c_{\text{top}} \left( \bigoplus_{j=1}^N R^1 \pi_* \mathcal{L}^{\otimes w_j} \right),
\]

for \( \phi_{b_i}, \ldots, \phi_{b_a} \) lying within the extended state space \( H_{\text{ext}} \).

\[\text{Proposition 3.2.} \quad \text{The above invariants vanish if one of the entries } \phi_{b_i}, \ldots, \phi_{b_a} \text{ does not belong to the narrow state space } H_{\text{nar}}. \text{ Otherwise } \langle \tau_{b_i}(\phi_{b_i}), \ldots, \tau_{b_a}(\phi_{b_a}) \rangle_{\text{FJRW}} \text{ equals } (\tau_{b_i}(\phi_{b_i}), \ldots, \tau_{b_a}(\phi_{b_a}))_{\text{FJRW}}.\]

\[\text{Proof.} \quad \text{The proof parallels the argument of [13, Lemma 4.1.1]. Let us compare } \mathcal{L}^{\otimes w_j} \text{ and } \mathcal{L}^{\otimes w_j} \text{ after push-forward via the morphism } \rho: \mathcal{C} \to \bar{\mathcal{C}} \text{ forgetting the stack-theoretic structure at the markings. We get}^{13}\]

\[\rho_*(\mathcal{L}^{\otimes w_j}) = \rho_*(\mathcal{L}^{\otimes w_j}) \otimes \mathcal{O} \left(- \sum_{i: (k_i+1)w_j \in d\mathbb{Z}} \bar{D}_i \right),\]

where \( \bar{D}_i \subset \bar{\mathcal{C}} \) is the divisor supported on the \( i \)th coarse marking. Therefore, if \((k_i+1)w_j \notin d\mathbb{Z}\) for all \( i \), we have \( R^1 \pi_*(\mathcal{L}^{\otimes w_j}) = R^1 \pi_*(\mathcal{L}^{\otimes w_j}) \). This shows the second claim above.

The vanishing condition in the statement holds when \((k_i+1)w_j \in d\mathbb{Z}\) for some \( 1 \leq i \leq n \) and some \( 1 \leq j \leq N \). This simply means that \( \mathcal{L}^{\otimes w_j}|_{\bar{D}_i} \) is pulled back from the coarse divisor \( \bar{D}_i \). On the other hand \( \mathcal{L}^{\otimes w_j} \) is a root of \( \omega_{\log}|_{\bar{D}_i} \cong \mathcal{O}_{\bar{D}_i} \) via the residue map. Therefore, \( c_1(\mathcal{L}^{\otimes w_j}|_{\bar{D}_i}) = 0 \) and hence \( c_1(\rho_*(\mathcal{L}^{\otimes w_j})|_{\bar{D}_i}) = 0 \) in the rational cohomology group.

Set \( T = \rho_* (\mathcal{L}^{\otimes w_j}) \). From (27), \( T(\bar{D}_i)|_{\bar{T}_i} = \rho_*(\mathcal{L}^{\otimes w_j})|_{\bar{T}_i} \) has vanishing first Chern class. Write the exact sequence

\[
0 \longrightarrow T \longrightarrow T(\bar{D}_i) \longrightarrow T(\bar{D}_i)|_{\bar{T}_i} \longrightarrow 0
\]

and the induced exact sequence of vector bundles

\[
0 \longrightarrow \pi_*(T(\bar{D}_i)|_{\bar{T}_i}) \longrightarrow R^1 \pi_*(T) \longrightarrow R^1 \pi_*(T(\bar{D}_i)) \longrightarrow 0.
\]

The vanishing \( c_{\text{top}}(R^1 \pi_*(\mathcal{L}^{\otimes w_j})) = c_{\text{top}}(R^1 \pi_*(T)) = 0 \) follows from \( c_1(T(\bar{D}_i)|_{\bar{T}_i}) = 0 \). Note that, in order to get (28), we need to show that \( T(\bar{D}_i) \) has only trivial sections on each fibre \( \bar{C} \) of \( \bar{\mathcal{C}} \). For \( a = d/w_j \), we find that \( T(\bar{D}_i) \otimes a \cong \omega(\bar{D}_i) - \sum_{i: k_i = a} \omega(\bar{D}_i) - \sum_{i: k_i \neq a} \omega(\bar{D}_i) \) which is a subsheaf of \( \omega(\bar{D}_i) \). It is easy to see that \( H^0(\bar{C}, \omega(\bar{D}_i)|_{\bar{T}}) = 0 \) for a genus zero curve \( \bar{C} \) by induction on the components (see [13]). Therefore \( H^0(\bar{C}, T(\bar{D}_i)|_{\bar{T}}) = 0 \). \( \Box \)

\[^{13}\text{To see this, use } \rho^* \rho_* \mathcal{E} = \mathcal{E} \otimes \mathcal{O}(-\sum_{i=1}^n d \text{age}_{\bar{D}_i}(\mathcal{E}) \bar{D}_i) \text{ for an invertible sheaf } \mathcal{E} \text{ on } \mathcal{C}.\]
3.2. Twisted FjRW theory and Givental’s formalism. — Let $K = \mathbb{C}[[s]]$ denote the completion of the polynomial ring $\mathbb{C}[s^j_k \mid 1 \leq j \leq N, k \geq 0]$ with respect to the additive valuation

$$v(s^j_k) = k + 1.$$ 

We define the ring $K[z, z^{-1}]$ of adically convergent power series in $z$ by

$$K[z, z^{-1}] = \left\{ \sum_{n \in \mathbb{Z}} a_n z^n \mid a_n \in K, v(a_n) \to \infty \text{ as } |n| \to \infty \right\}.$$ 

Define $K[z]$ (resp. $K[z^{-1}]$) to be the subring of $K[z, z^{-1}]$ consisting of non-negative (resp. non-positive) power series in $z$. We introduce a symmetric non-degenerate pairing $(\cdot, \cdot)_s$ on $H_{\text{ext}} \otimes K$ taking values in $K$:

$$(\phi_h, \phi_k)_s = \frac{1}{d} \left( \prod_{j=0}^{h+k} \exp(-s^j_0) \right) \delta_{h+k, d-2},$$

where $\delta_{h+k, d-2}$ is 1 if $h + k \equiv d - 2 \mod d$ and 0 otherwise. Through the entire section we adopt the convention that the index is reduced modulo $d$ to the suitable range $\{0, \ldots, d - 1\}$.

**Definition 3.3 (Givental’s symplectic space).** — Givental’s symplectic space is the space

$$\mathcal{H} := H_{\text{ext}} \otimes K[z, z^{-1}]$$

equipped with the symplectic form

$$\Omega^s(f_1, f_2) = \text{Res}_{z=0}(f_1(-z), f_2(z)) \, dz.$$ 

The space $\mathcal{H}$ has a standard polarization $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-$, where $\mathcal{H}_+ = H_{\text{ext}} \otimes K[z]$, $\mathcal{H}_- = H_{\text{ext}} \otimes z^{-1} K[z^{-1}]$ are isotropic subspaces of $\mathcal{H}$. This polarization allows us to identify $\mathcal{H}$ with the total space of the cotangent bundle of $\mathcal{H}_+$.

For the basis $\{\phi_i\}_{i=0}^{d-1}$ of $H_{\text{ext}}$, we write $g^s_{\phi_k}$ for $(\phi_h, \phi_k)_s$ and $g^{ph}_{\phi_k}$ for the coefficients of the inverse matrix. A general point of $\mathcal{H}$ can be written as $\mathbf{q} + \mathbf{p}$ with

$$(30) \quad \mathbf{q} = \sum_{b \geq 0} \sum_{k=0}^{d-1} q^k_b \phi_k z^b \in \mathcal{H}_+, \quad \mathbf{p} = \sum_{b \geq 0} \sum_{h,k=0}^{d-1} p_{b,h} g^{ph}_{\phi_k} (-z)^{1+b} \in \mathcal{H}_-.$$ 

Here $\{q^k_b, p_{b,k} \mid b \geq 0, 0 \leq k \leq d - 1\}$ can be regarded as Darboux co-ordinates on $\mathcal{H}$. Following established practice, we denote the coordinates in Givental formalism by $q^k_b$, with the label $b$ corresponding to gravitational descendants and with the label $k$ corresponding to state space entries. We indicate explicitly where the superscript “$k$” index is meant instead as an exponent of a power, indeed this turns out to be useful in some very special cases.
Definition 3.4 (Twisted FJRW theory cf. [21]). Consider the universal characteristic class of the extended obstruction $K$-class:

\[
e(s) = \exp \left( \sum_{1 \leq j \leq N} \sum_{l=0}^{\infty} s_l^{(j)} \text{ch}_1(R \tau_s \tilde{L}^\otimes w) \right) \in H^*(\text{Spin}^d_0(k_1, \ldots, k_n); \mathbb{C}) \otimes K
\]

and define the twisted FJRW invariants as

\[
\langle \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \rangle_{s,0} = \int_{[\text{Spin}^d_0(k_1, \ldots, k_n)]} \left( \prod_{i=1}^{n} \psi_i^{b_i} \right) \cup e(s).
\]

The twisted FJRW invariants are encoded in the generating function

\[
F_0^s = \sum_{b_1, \ldots, b_n \geq 0 \atop 0 \leq k_1, \ldots, k_n \leq d-1} \langle \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \rangle_{0,n} \frac{t_{k_1}^{b_1} \cdots t_{k_n}^{b_n}}{n!}.
\]

This is a formal power series in infinitely many variables \( \{ t_k^b \mid 0 \leq k \leq d-1, b \geq 0 \} \). Again, the superscript \( k \) of \( t_k^b \) means an index, not an exponent of a power.

The twisted FJRW invariants here are a generalization of the extended invariants (26).

Definition 3.5 (Givental’s Lagrangian submanifold). We relate the variables \( \{ t_k^b \} \) of \( F_0^s \) and the co-ordinates \( \{ q_k^b \} \) on \( \mathcal{H}_+ \) by the following dilaton shift:

\[
q_k^b = -\delta_0^k \delta_0^k + t_k^b.
\]

Then \( F_0^s \) can be regarded as a function defined on a formal neighbourhood of \( -z\phi_0 \in \mathcal{H}_+ \). The graph of \( dF_0^s \) defines a Lagrangian submanifold of \( (\mathcal{H}, \Omega^s) \):

\[
L^s := \left\{ q + p \in \mathcal{H} \mid p_{b,k} = \frac{\partial F_0^s}{\partial q_k^b}, b \geq 0, 0 \leq k \leq d-1 \right\}.
\]

The submanifold \( L^s \) can be defined as a formal scheme over \( K \). See [17, Appendix B].

3.2.1. The untwisted theory. Consider the case where \( s_l^{(j)} = 0 \) for all \( 1 \leq j \leq N \) and \( l \geq 0 \). Then \( e(s) = 1 \) and the associated correlators give the so called untwisted invariants...
\[
\left\langle \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \right\rangle_{0,n}^{\text{un}}
= \int_{[\text{Spin}_{d,n}^0(k_1, \ldots, k_n)]} \prod_{i=1}^{n} \psi_i^{b_i}
= \begin{cases}
\frac{1}{d} \frac{(\sum_{i=1}^{n} k_i)!}{b_1! \cdots b_n!} & \text{if } n - 3 = \sum_{i=1}^{n} b_i \text{ and } 2 + \sum_{i=1}^{n} k_i \in d\mathbb{Z}, \\
0 & \text{otherwise}.
\end{cases}
\]

One can show this by using the String Equation (since one of \( b_i \) has to be zero). The same Hodge integral over \( M_{0,n} \) is given in [24, Equation (2)]; the new factor \( 1/d \) here comes from the fact that \( \text{Spin}_{d,n}^0(k_1, \ldots, k_n) \) has \( \mu_d \) as the generic stabilizer and that \( [\text{Spin}_{d,n}^0(k_1, \ldots, k_n)] = \frac{1}{d}[M_{0,n}] \). The generating function \( F_0^{\text{un}} \) of untwisted invariants are defined similarly to (33). The pairing \( (\cdot, \cdot)_s \) and the symplectic form \( \Omega^s \) specialize to
\[
(\phi_k, \phi_h)_s = \frac{1}{d} \delta_{d-2,k+h}, \quad \Omega^s(f_1, f_2) = \text{Res}_{z=0} (f_1(-z), f_2(z))\bigg|_{0}^{d}.
\]

The Lagrangian submanifold \( L^{\text{un}} \subset (\mathcal{H}, \Omega^{\text{un}}) \) is defined as the graph of \( dF_0^{\text{un}} \) as in (34). (Here one should use as Darboux co-ordinates those given by the untwisted pairing \( g_{s,k}^{\text{un}} = (\phi_k, \phi_h)_\text{un} \) instead of \( g_{s,k}^{\text{tw}} \), cf. (30).)

Since the untwisted invariants are the usual intersection numbers on \( \overline{M}_{0,n} \), the generating function \( F_0^{\text{un}} \) satisfy the well known tautological equations: String Equation (SE), Dilaton Equation (DE) and Topological Recursion Relations (TRR). Givental [29] showed that these three equations for a genus zero potential \( F_0 \) are equivalent to the following geometric properties for the graph \( L \) of the differential \( dF_0^\text{un} \):

- \( L \) is a cone in \( \mathcal{H} \) with vertex at the origin \( p = q = 0 \) (with the dilaton shift \( q_a^k = t_a^k - \delta_a^1 \delta^k_0 \) understood);
- The tangent space \( T \) to \( L \) at any point on \( L \) satisfies \( zT = L \cap T \); Moreover the tangent space to \( L \) at any point in \( zT \subset L \) equals \( T \).

We refer to these properties as Givental’s geometric properties for \( L \). In particular, \( L^{\text{un}} \) satisfies Givental’s geometric properties.

**3.2.2. The twisted theory.** — The Lagrangian submanifold \( L^s \) was determined by Chiodo-Zvonkine [16]. Define a linear symplectic transformation \( \Delta : (\mathcal{H}, \Omega^{\text{un}}) \to (\mathcal{H}, \Omega^s) \) by

\[
\Delta = \bigoplus_{i=0}^{d-1} \exp \left( \sum_{j=1}^{N} \sum_{l \geq 0} s_j^i B_{l+1}((iq_j) + q_j) \frac{z^j}{(l+1)!} \right)
\]

where \( B_n(x) \) is the Bernoulli polynomial defined by \( \sum_{n=0}^{\infty} B_n(x) z^n/n! = ze^{xz}/(e^z - 1) \).
Theorem 3.6 (Chiodo-Zvonkine [16]). — We have $\mathcal{L}^s = \Delta(\mathcal{L}^\text{un})$.

Because Givental’s geometric properties are preserved by a linear symplectic transformation, the generating function $\Phi^s_0$ of twisted FJRW invariants satisfies $\text{SE}$, $\text{DE}$ and $\text{TRR}$.

The adaptation of [16] to our context was explained in [13, Proposition 4.1.5]; we omit the details.

3.3. Family of elements on the Lagrangian cone. — The twisted $J$-function is a family of elements lying on $\mathcal{L}^s$ parametrized by $t = \sum_{k=0}^{d-1} t^k \phi_k \in H_{\text{ext}} \otimes K$:

$$J^s(t, -z) = -z \phi_0 + t$$

$$+ \sum_{a=2}^{\infty} \sum_{b=0}^{\infty} \sum_{0 \leq j \leq d-1} \frac{1}{n!} (t, \ldots, t, \tau_0(\phi_k))_{a+b}(s \phi_k (-z))^{b+1}.$$  

Here $J^s(t, -z) \in \mathcal{H}$ is characterized as a unique point lying on $\mathcal{L}^s$ with the property:

$$J^s(t, -z) = -\phi_0 z + t + O(z^{-1}).$$

It is known [29] that the $J$-function reconstructs the cone $\mathcal{L}^s$ itself via Givental’s geometric properties. Here we will find another explicit family of elements (I-function) on $\mathcal{L}^s$.

The $J$-function $J^\text{un} \in \mathcal{L}^\text{un}$ of the untwisted theory (Section 3.2.1) is the specialization of (37) at $s = 0$. Using (35), we calculate

$$J^\text{un}(t, -z) = \sum_{k = (k_0, \ldots, k_{d-1}) \in \mathbb{Z}_{\geq 0}^{d-1}} J^\text{un}_k(t, -z),$$

$$J^\text{un}_k(t, -z) = \frac{1}{(-z)^{|k|-1}} \frac{(t^k)^{k_0} \ldots (t^{d-1})^{k_{d-1}}}{k_0! \ldots k_{d-1}!} \phi_{h(k)},$$

with $|k| = \sum_{i=0}^{d-1} k_i$ and $h(k) = \sum_{i=0}^{d-1} i k_i$. Here $(t^k)^i$ means the $k_i$-th power of the variable $t^i$. Introduce the modification factor $M_k(z)$ by

$$M_k(z) = \prod_{j=1}^{N} \exp \left( - \sum_{0 \leq m < [q_j h(k)]} s^{(j)}_m (-(q_j + [q_j h(k)] + m) z) \right),$$

where $s^{(j)}_n = \sum_{a \geq 0} s_n^{(j)} x^n/n!$ and define the twisted I-function by

$$I^s(t, z) = \sum_{k_0, \ldots, k_{d-1} \geq 0} M_k(z) J^\text{un}_k(t, z).$$

Using Theorem 3.6, we get the following statement.

Theorem 3.7. — The family $t \mapsto I^s(t, -z)$ of elements of $\mathcal{H}$ lies on $\mathcal{L}^s$. 
Proof. — The discussion here is parallel to [17, Theorem 4.8] and [13, Theorem 4.1.6]. We give a sketch of the proof and leave the details to the reader. Introduce a function

\[ G^{(j)}_y(x, z) = \sum_{m, l \geq 0} s^{(j)}_{l+m-1} \frac{B_m(y)}{m!} \frac{x^l}{l!} z^{m-1}, \quad j = 1, \ldots, N \]

with \( s^{(j)}_{-1} = 0 \). Set \( D := \sum_{k=0}^{d-1} k t^k (\partial / \partial t^k) \). Givental’s geometric properties for the cone \( L^m \) yield the following fact (see [17, Equation (14)] and [13, Lemma 4.1.10]):

Lemma 3.8. — The family \( t \mapsto \exp(- \sum_{j=1}^N G^{(j)}_0 (zq_j D + zq_j, z)) J^m(t, -z) \) lies on \( L^m \).

The conclusion of Theorem 3.7 follows from Theorem 3.6: we apply the symplectic transformation \( \Delta : (\mathcal{H}, \Omega^m) \to (\mathcal{H}, \Omega^*) \) in (36) to the family in Lemma 3.8. Note that we have

\[ \Delta = \bigoplus_{i=0}^{d-1} \exp \left( \sum_{j=1}^N G^{(j)}_y (0, z) \right) = \bigoplus_{i=0}^{d-1} \exp \left( \sum_{j=1}^N G^{(j)}_0 ((i q_j) + q_j) z, z \right) \]

where we used \( G^{(j)}_y(x, z) = G^{(j)}_0(x + y z, z) \) in the second equality. Using the identity

\[ G^{(j)}_0(x + z, z) = G^{(j)}_0(x, z) + s^{(j)}(x) \]

we can easily check that

\[ \Pi^*(t, -z) = \Delta \exp \left( - \sum_{j=1}^N G^{(j)}_0 (zq_j D + zq_j, z) \right) J^m(t, -z). \]

Theorem 3.6 and Lemma 3.8 show that \( \Pi^*(t, -z) \) is on the cone \( L^* \). \( \square \)

3.4. The twist by the equivariant Euler class. — Let \( T = (\mathbb{C}^\times)^N \) act on the extended obstruction bundle \( \bigoplus_{j=1}^N R^1 \pi_* (\widetilde{L}^\otimes w_j) \) diagonally by scaling the fibres and trivially on the base \( \text{Spin}^d_{0, n}(k_1, \ldots, k_n) \). Let \( \lambda_1, \ldots, \lambda_N \in H^2_T(\text{pt}) \) denote the equivariant parameters. Then the \( T \)-equivariant Euler class \( e_T \) of the extended obstruction bundle is given by

\[ e_T \left( \bigoplus_{j=1}^N R^1 \pi_* (\widetilde{L}^\otimes w_j) \right) = \prod_{j=1}^N \sum_{l=0}^{\gamma_j - 1} \lambda_j^{\gamma_j - l} \epsilon_l \left( R^1 \pi_* (\widetilde{L}^\otimes w_j) \right) \]

with \( \gamma_j = \text{rank}(R^1 \pi_* (\widetilde{L}^\otimes w_j)) \). This class can be obtained from the universal class \( e(s) \) (31) by the substitution:

\[ s^{(j)}_l = \begin{cases} -\log \lambda_j & l = 0; \\ (l - 1)! (\gamma_j - l)^{-l} & l > 0. \end{cases} \]
Note that this specialization yields \( \exp(-s^{(j)}(x)) = x + \lambda_j \), where \( s^{(j)}(x) = \sum_{n=0}^{\infty} \frac{s^{(j)}(x)^n}{n!} \) as before. With this choice of parameters, we obtain

- The \( e_T\)-twisted pairing as the specialization of (29):
  \[
  (\phi_b, \phi_k)_{\text{tw}} := \frac{1}{d} \left( \prod_{j:(q,k) = 1} \lambda_j \right) \delta_{d-2,k+b};
  \]

- The \( e_T\)-twisted FJRW invariants \( \{ \tau_{h_1}(\phi_{b_1}), \ldots, \tau_{b_n}(\phi_{b_n}) \}_{\text{tw}} \) as the specializations of (32);

- The \( e_T\)-twisted \( J \)-function \( \Gamma^{(w)}(t, -z; \lambda) \) as the specialization of (37);

- The \( e_T\)-twisted I-function \( \Gamma^{(w)}(u, z; \lambda) \) as the specialization of \( uI^*(u \phi_1, z) \) (see (39)):

\[
\Gamma^{(w)}(u, z; \lambda) := z \sum_{k=0}^{\infty} u^k \prod_{j=1}^{N} \prod_{0 < k < d, (b) = (q,k)} \left( \lambda_j - b \right) \frac{\prod_{0 < k < d} (-bz)}{\prod_{0 < k < d} (-bz)} \phi_{k-1}.
\]  
Notice that the denominator \( \prod_{0 < k < d} (-bz) \) here is nothing but \( (k-1)!(-z)^{k-1} \), but we prefer this expression in view of our parallel treatment of I-functions in Section 4.3.1.

Remark 3.9. — Note that the specialization \( s^{(j)}_0 = -\log \lambda_j \) does not make sense for every element in the ground ring \( K \). For this reason, we do not try to define the \( e_T\)-twisted Lagrangian cone. The specializations of the I- and J-functions, however, still make sense as elements of \( H_{\text{ext}} \otimes \mathbb{C}[z, z^{-1}][\lambda_1^\pm, \ldots, \lambda_n^\pm][e^0, \ldots, e^{d-1}] \).

The \( e_T\)-twisted I-function has the following \( z \)-asymptotics

\[
\Gamma^{(w)}(u, z; \lambda) = zF(u)\phi_0 + G(u; \lambda) + O(z^{-1})
\]
where \( F \in \mathbb{C}[[u]] \) is a scalar-valued and \( G \in H_{\text{ext}}^{<2} \otimes \mathbb{C}[\lambda_1, \ldots, \lambda_N][u] \) is an \( H_{\text{ext}}^{<2} \)-valued power series (where \( H_{\text{ext}}^{\leq 2} \) denotes the degree \( \leq 2 \) part of \( H_{\text{ext}} \)):

\[
F(u) = \sum_{k \geq 1 : d \equiv 1 (d)} u^k \frac{\prod_{j=1}^{N} (kq_j - 1)_{[kq_j]_{d-1}}}{(k-1)!},
\]

\[
G(u; \lambda) = -\sum_{k \geq 2 : \sum_{j=1}^{N} (q_j(k-1)) = 1} u^k \frac{\prod_{j=1}^{N} (kq_j - 1)_{[kq_j]_{d-1}}}{(k-1)!} \phi_{k-1}
\]

\[
-\sum_{i=1}^{N} \lambda_i \sum_{k \geq d+1 : d \equiv 1 (d)} u^k \left( \sum_{0 < a < kq_j, (a) = \{kq_j\}} \frac{1}{a} \right) \prod_{j=1}^{N} \frac{(kq_j - 1)_{[kq_j]_{d-1}}}{(k-1)!} \phi_0,
\]
where \((a)_n = a(a-1)\cdots(a-n+1) = \Gamma(a+1)/\Gamma(a-n+1)\) denotes the falling factorial. We define the \(FJRW\) mirror map to be the \(H^{\leq 2}_{\text{ext}}\)-valued function:

\[
\zeta(u; \lambda) = \frac{G(u; \lambda)}{F(u)} = -u \phi_1 + O(u^2) \in H^{\leq 2}_{\text{ext}} \otimes \mathbb{C}[\lambda_1, \ldots, \lambda_N][[u]].
\]

We now state the mirror theorem for the \(e_1\)-twisted \(FJRW\) theory.

**Theorem 3.10.** — We have \(J^{tw}(\zeta(u; \lambda), z; \lambda) = I^{tw}(u, z; \lambda)/F(u)\) for the function \(F(u)\) in (42).

**Proof.** — Because of the problem we discussed in Remark 3.9, we use another specialization \(\psi_t^{(j)} = \bar{s}_t^{(j)}\), where

\[
\bar{s}_t^{(j)} := \begin{cases} 0 & l = 0; \\ (l - 1)!(-\lambda_j)^{-l} & l \geq 0. \end{cases}
\]

This specialization yields \(\exp(-\bar{s}_t^{(j)}(x)) = 1 + x/\lambda_j\). It defines a well-defined homomorphism \(K \to \mathbb{C}[[\lambda^{-1}]] := \mathbb{C}[[\lambda_1^{-1}, \ldots, \lambda_N^{-1}]]\) and the characteristic class:

\[
e(\bar{s}) = \left(\prod_{j=1}^{N} \lambda_j^{-\bar{s}_t^{(j)}}\right) e_1 \left(\bigoplus_{j=1}^{N} R^1 \pi_* \left(\bar{L}^{\otimes \bar{s}_t^{(j)}}\right)\right).
\]

The Lagrangian cone \(L^{\bar{s}}\) can be defined as a formal scheme over \(\mathbb{C}[[\lambda^{-1}]]\) and \(I^{\bar{s}}(t, -z)\) is lying on \(L^{\bar{s}}\) by Theorem 3.7. After some computation, we find that \(J^{tw}\) and \(I^{tw}\) are related to \(J^{\bar{s}}\) and \(I^{\bar{s}}\) as

\[
J^{tw}(t, z; \lambda) = R(\lambda)J^{\bar{s}}(R(\lambda)^{-1}t, z; \lambda)
\]

\[
I^{tw}(u, z; \lambda) = R(\lambda)uI^{\bar{s}}(-uR(\lambda)^{-1}\phi_1, z)
\]

where \(R(\lambda): H_{\text{ext}} \to H_{\text{ext}}\) is defined by \(R(\lambda)\phi_h = (\prod_{j=1}^{N} \lambda_j^{-\langle h, \lambda_j \rangle})\phi_h\). It is also easy to check that \(I^{\bar{s}}(-u\phi_1, z)\) has the asymptotics

\[
I^{\bar{s}}(-u\phi_1, z) = zF(u; \lambda)\phi_0 + \bar{G}(u; \lambda) + O(z^{-1})
\]

for a scalar valued function \(\bar{F} = 1 + O(u)\) and an \(H^{\leq 2}_{\text{ext}}\)-valued function \(\bar{G}\). Here the functions \(F, G\) appearing in (42) are related to \(\bar{F}, \bar{G}\) as

\[
F(u) = u\bar{F}(\lambda^g u; \lambda), \quad G(u; \lambda) = uR(\lambda)\bar{G}(\lambda^g u; \lambda)
\]

with \(\lambda^g = \prod_{j=1}^{N} \lambda_j^{\bar{s}_t^{(j)}}\). Because \(L^{\bar{s}}\) is a cone and \(I^{\bar{s}}(-u\phi_1, -z)\) is on \(L^{\bar{s}}\), we have (we regard \(I^{\bar{s}}(-u\phi_1, -z)\) as a \(\mathbb{C}[[\lambda^{-1}]][[u]]\)-valued point on \(L^{\bar{s}}\) and apply [17, Proposition B2]):

\[
\frac{1}{F(u; \lambda)} I^{\bar{s}}(-u\phi_1, -z) = -z\phi_0 + \frac{\bar{G}(u; \lambda)}{F(u; \lambda)} + O(z^{-1}) \in L^{\bar{s}}.
\]
By the characterization (38) of the J-function, this coincides with \( \overline{\mathcal{G}(u; \lambda)} / \overline{\mathcal{F}(u; \lambda)} - z \).

The conclusion follows from this and the relations (43), (44).

3.5. The \( e_{\mathbb{C}^*} \)-twisted quantum connection. — Here we discuss the \( e_{\mathbb{C}^*} \)-twisted quantum cohomology for both of FJRW and GW theory. We show that the non-equivariant limit \( \lambda \to 0 \) of the \( e_{\mathbb{C}^*} \)-twisted quantum product exists and reduces to the original one (17) restricted to the narrow/ambient part. In the rest of the paper, we only consider the \( e_{\mathbb{C}^*} \)-twisted theory as a twisted theory. Thus the word “\( e_{\mathbb{C}^*} \)-twisted” is sometimes abbreviated as “twisted”.

3.5.1. A brief summary of the \( e_{\mathbb{C}^*} \)-twisted theory. — We mean by the \( e_{\mathbb{C}^*} \)-twisted FJRW theory the \( e_T \)-twisted FJRW theory (Section 3.4) with the equivariant parameters \( \lambda_1, \ldots, \lambda_N \) specialized to the following values:

\[
\lambda_i = -q_i \lambda, \quad i = 1, \ldots, N.
\]

The \( e_{\mathbb{C}^*} \)-twisted pairing \( (\phi_+, \phi_-)_t \) and FJRW invariants \( \langle \tau_{b_1}(\phi_{b_1}), \ldots, \tau_{b_n}(\phi_{b_n}) \rangle_{0,n}^{\text{FJRW},\text{tw}} \) take values in \( \mathbb{C}[\lambda] \). (We have put the superscript “FJRW” to distinguish them from GW invariants.)

The \( e_{\mathbb{C}^*} \)-twisted GW theory [17, 21, 66] for \( \mathcal{P}(w) \) is defined on the state space \( \text{HCR}(\mathcal{P}(w)) = \text{H}(\mathcal{I}\mathcal{P}(w)) \). We consider the twist by the line bundle \( O(d) \). Let \( \mathcal{P}(w)_{0,n,\beta} \) denote the moduli stack of \( n \)-pointed stable maps to \( \mathcal{P}(w) \) of genus 0 and degree \( \beta \). Let \( \pi : \mathcal{C}_{0,n,\beta} \to \mathcal{P}(w)_{0,n,\beta} \) be the universal curve and let \( f : \mathcal{C}_{0,n,\beta} \to \mathcal{P}(w) \) be the universal stable map:

\[
\begin{array}{c}
\mathcal{C}_{0,n,\beta} \quad \xrightarrow{f} \quad \mathcal{P}(w) \\
\downarrow \pi \\
\mathcal{P}(w)_{0,n,\beta}
\end{array}
\]

It can be shown that \( R\pi_* f^* O(d) \) is represented by a vector bundle (see [20]; this is because \( O(d) \) is ample and \( O(d) \) is pulled back from the coarse moduli \( |\mathcal{P}(w)| \)). The \( e_{\mathbb{C}^*} \)-twisted GW invariants of \( \mathcal{P}(w) \) are defined to be

\[
\left\{ \tau_{b_1}(\alpha_1), \ldots, \tau_{b_n}(\alpha_n) \right\}_{0,n,\beta}^{\text{GW},\text{tw}} = \int_{[\mathcal{P}(w)_{0,n,\beta}]} \prod_{i=1}^n \text{ev}_i^*(\alpha_i) \psi_i^{b_i} \cup e_{\mathbb{C}^*} \left( R\pi_* f^* O(d) \right),
\]

where \( \alpha_1, \ldots, \alpha_n \in \text{HCR}(\mathcal{P}(w)) \) and \( \mathbb{C}^* \) acts on \( R\pi_* f^* O(d) \) by scaling the fibre. This is an element of \( \text{H}_{\mathbb{C}^*}(\text{pt}) = \mathbb{C}[\lambda] \). We endow \( \text{H}_{\mathbb{C}^*}(\mathcal{P}(w)) \) with the following twisted pairing

\[
(\alpha_1, \alpha_2)_{\text{tw}} = \int_{\mathcal{I}\mathcal{P}(w)} \alpha_1 \cup \alpha_2 \cup e_{\mathbb{C}^*} \left( \text{pr}^* O(d) \right)
\]

where \( \text{pr} : \mathcal{I}\mathcal{P}(w) \to \mathcal{P}(w) \) is the natural projection.
3.5.2. Twisted quantum product. — We denote by $\overline{H}$ the state space of the twisted theory:

$$
\overline{H} := \begin{cases} 
H_{\text{ext}} & \text{for FJRW theory;} \\
H_{\text{CR}}(\mathcal{P}(w)) & \text{for GW theory.}
\end{cases}
$$

Both state spaces are of dimension $d$. The same procedure as Section 2.3 defines the twisted quantum cohomology. The $e_{\mathbb{C}^\times}$-twisted quantum product $\bullet_{\text{tw}}$ on $\overline{H}$ is defined by the formula (17) with the correlators $\langle \cdots \rangle_{0,n}$ replaced by the $e_{\mathbb{C}^\times}$-twisted invariants and $(g_{\rho})$ replaced by the $e_{\mathbb{C}^\times}$-twisted pairing. Because the divisor equation holds also for the twisted GW theory, we can consider the specialization $Q = 1$ for $\bullet_{\text{tw}}$ (see Section 2.3.2). In GW theory, we shall denote by $\bullet_{\text{tw}}$ the twisted quantum product with $Q$ already specialized to 1.

Let $H'$ denote the narrow/ambient part (21) of the state space. Let $pr$ denote the natural projection

$$
pr: \overline{H} \rightarrow H'.
$$

Let $T_0, \ldots, T_{d-1}$ be a homogeneous basis of $\overline{H}$ such that $T_0$ is the identity ($T_0 = \phi_0$ in FJRW theory$^{14}$ and $T_0 = 1_0$ in GW theory). In the case of GW theory, we take $T_1 = p = c_1(\mathcal{O}(1))$. Let $t^0, \ldots, t^{d-1}$ denote the linear co-ordinate on $\overline{H}$ dual to the basis $T_0, \ldots, T_{d-1}$.

**Proposition 3.11.** — The $e_{\mathbb{C}^\times}$-twisted quantum products are regular at $\lambda = 0$, i.e.

$$
T_i \bullet_{\text{tw}} T_j \in \begin{cases} 
\overline{H} \otimes \mathbb{C}[\lambda][[t^0, \ldots, t^{d-1}]] & \text{for FJRW theory;} \\
\overline{H} \otimes \mathbb{C}[\lambda][[t^0, e^1/w, t^2, \ldots, t^{d-1}]] & \text{for Gromov-Witten theory.}
\end{cases}
$$

Here $w$ is the least common multiple of $w_1, \ldots, w_n$ (see Section 2.3.2). Moreover we have

$$
\lim_{\lambda \to 0} pr(T_i \bullet_{t} T_j) = pr(T_i) \bullet_{pt(t)} pr(T_j)
$$

where the product in the right-hand side is the ordinary quantum product on the narrow/ambient part in Section 2.4 and the subscripts $t \in \overline{H}$, $pt(t) \in H'$ denote the parameter of the product.

**Proof.** — This was proved in [43, Corollary 2.5] for GW theory, so we only discuss the case of the FJRW theory. The $e_{\mathbb{C}^\times}$-twisted FJRW quantum product can be written as

$$
\phi_i \bullet_{\text{tw}} \phi_j = \sum_{k=0}^{d-1} \sum_{n \geq 0} \frac{1}{n!} \langle \phi_i, \phi_j, \phi_k, t, \ldots, t \rangle_{0,n+3}^{\text{FJRW,}\text{tw}} \left( d \prod_{j:(k+1)g_j = 0} \lambda_j^{-1} \right) \phi_{d-2-k}
$$

$^{14}$ The element $\phi_0$ is the identity in the twisted FJRW theory because of the string equation (see Section 3.2.2).
with $\lambda_j = -q_j \lambda$. To see that this expression is regular at $\lambda = 0$, it suffices to show that
\[
\{ \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \}_{0,n}^{\text{FJRWW},t} \in \left( \prod_{j;(k_i+1)q_j=0} \lambda_j \right) \mathbb{C}[\lambda_1, \ldots, \lambda_N],
\]
\[1 \leq \forall i \leq n.
\]
This happens because $e(\mathcal{R}_1 \pi^* \tilde{L} \otimes w_j)$ is divisible by $\lambda_j$ as soon as $d$ divides $(k_i+1)w_j$. This follows from the fact that $\mathcal{R}_1 \pi^* \tilde{L} \otimes w_j$ contains the sub-line bundle $\pi^* (T(D_i)|_{D_i})$ whose equivariant 1st Chern class is $\lambda_j$. See (28).

Using Proposition 3.2 and the fact that the $e_{\mathbb{C}^\times}$-twisted invariant equals the extended invariant (26) in the non-equivariant limit $\lambda \rightarrow 0$, we have
\[
\lim_{\lambda \rightarrow 0} \{ \tau_{b_1}(\phi_{k_1}), \ldots, \tau_{b_n}(\phi_{k_n}) \}_{0,n}^{\text{FJRWW},t} = \{ \tau_{b_1}(\text{pr}(\phi_{k_1})), \ldots, \tau_{b_n}(\text{pr}(\phi_{k_n})) \}_{0,n}^{\text{FJRWW}}.
\]

The equality (45) follows easily from this. □

3.5.3. Twisted quantum connection and the fundamental solution. — The $e_{\mathbb{C}^\times}$-twisted quantum connection is defined similarly to (18):
\[
\nabla_i^{\text{tw}} = \frac{\partial}{\partial t_i} + \frac{1}{z} T_i \cdot \text{tw}.
\]

By Proposition 3.11, $e_{\mathbb{C}^\times}$-twisted quantum connection is regular at $\lambda = 0$. The non-equivariant limit is called the $e$-twisted quantum connection. In contrast with the untwisted case, the connection $\nabla^{\text{tw}}$ cannot be extended in the $z$-direction since the variable $\lambda$ has a degree.

Let $L^{\text{tw}}(t, z; \lambda)$ denote the canonical fundamental solution of the connection $\nabla^{\text{tw}}$ defined by the same formula (20) with $\langle \cdots \rangle_{0,n}, g_0$ there replaced with the twisted counterparts. This satisfies (part of) the properties in Proposition 2.7:

**Proposition 3.12.** — For $\alpha, \alpha_1, \alpha_2 \in \overline{\mathbb{H}}$, we have
\[
\nabla_i^{\text{tw}} L^{\text{tw}}(t, z; \lambda) \alpha = 0, \quad (L^{\text{tw}}(t, -z; \lambda) \alpha_1, L^{\text{tw}}(t, z; \lambda) \alpha_2)^{\text{tw}} = (\alpha_1, \alpha_2)^{\text{tw}}.
\]

In particular, the connection $\nabla^{\text{tw}}$ is flat, i.e. $[\nabla_i^{\text{tw}}, \nabla_j^{\text{tw}}] = 0$ and that the pairing $(\cdot, \cdot)^{\text{tw}}$ is $\nabla^{\text{tw}}$-flat (see (19) for a precise meaning of the flatness of the pairing).

**Proof.** — The outline of the proof is the same as Proposition 2.7. It suffices to show that the twisted theory satisfies TRR, but this follows from Givental’s geometric properties (see Section 3.2.2). The same discussion as in [43, Proposition 2.1] shows the statement for the pairing. □
Remark 3.13. — The fact that $\nabla^{tw}$ is flat implies that $\bullet^{tw}$ is associative. (The commutativity of $\bullet^{tw}$ is clear from the definition.)

Because $L^{tw}$ satisfies the differential equation regular at $\lambda = 0$, it follows that $L^{tw}$ is also regular at $\lambda = 0$ (see also [43, Proposition 2.4]).

\[
L^{tw}(t, z; \lambda) \in \begin{cases} 
\text{End}(H) \otimes \mathbb{C}[\lambda][[t^0, \ldots, t^{d-1}]][[z^{-1}]] \\
\text{for FJRW theory;}
\end{cases}
\]

for FJRW theory;

\[
\text{End}(H) \otimes \mathbb{C}[\lambda][[t^0, e^{t^1/w}, t^2, \ldots, t^{d-1}]][[t^1]][[z^{-1}]]
\]

for GW theory.

Let $L(t, z)$ denote the fundamental solution (20) in the original FJRW/GW theories. (For GW theory, we specialize $Q$ to 1.) We have the following:

Proposition 3.14.

\[
\lim_{\lambda \to 0} \text{pr}(L^{tw}(t, z; \lambda) \alpha) = L(\text{pr}(t), z) \text{pr}(\alpha).
\]

Proof. — It was shown in [43, Proposition 3.24] for GW theory. For FJRW theory, the equality follows easily from (46).

3.5.4. Twisted J-function. — Recall from Section 3.3 that the J-function (37) is a special family of elements lying on the Givental Lagrangian cone. The $\epsilon_{\mathbf{e}^*}$-twisted J-function is defined similarly:

\[
J^{tw}(t, z) = zT_0 + t + \sum_{n \geq 0} \sum_{b \geq 0} \sum_{i,j=0}^i \frac{1}{n!} z^{b+1} t, \ldots, t, \tau_b(T_i)^{tw}_{0,n+1} g^{tw}_{ij} T_j,
\]

where $(g_{ij}^{tw})$ denotes the inverse of the twisted pairing matrix $g^{tw}_{ij} = (T_i, T_j)_{tw}$. (In the case of GW theory, as in (16), we also take the summation over curve classes $\beta$ (see [17, Equation (8)]). Then we specialize it to $Q = 1$ using the divisor equation.) The following relation of $L^{tw}$ and $J^{tw}$ is a key to understand the role of the J-function in the quantum D-module.

Proposition 3.15. — $L^{tw}(t, z; \lambda)J^{tw}(t, z; \lambda) = zT_0$.

Proof. — By Proposition 3.12, we have $L^{tw}(t, z; \lambda)^{-1} = L^{tw}(t, -z; \lambda)^*$ where $*$ denotes the adjoint with respect to the twisted pairing. Thus we have

\[
(T_i, L_{tw}(t, z; \lambda)^{-1} zT_0)_{tw} = (T_i, L^{tw}(t, -z; \lambda)^* zT_0)_{tw}
\]

\[
= (L^{tw}(t, -z; \lambda) T_i, T_0)_{tw} = (T_i, J^{tw}(t, z; \lambda))_{tw},
\]

where the last equality follows directly from the definitions (20), (47) of $L^{tw}$ and $J^{tw}$ and the string equation. The conclusion follows.
4. Orlov equivalence matches Mellin-Barnes analytic continuation

4.1. Matrix factorizations. — Matrix factorizations were originally introduced by Eisenbud [23] for the study of maximal Cohen-Macaulay modules. Recently Kontsevich proposed that they form the category of B-branes in the Landau-Ginzburg model. References are made to [22, 34, 37, 52, 53, 56, 68]. The paper [22] contains a nice introduction to the subject.

We introduce the differential graded (dg) category of graded matrix factorizations of a degree-\(d\) weighted homogeneous polynomial \(W \in \mathbb{C}[x_1, \ldots, x_N]\) from the introduction Section 1.1. Set \(R := \mathbb{C}[x_1, \ldots, x_N]\). Notice that \(R\) is a \(\mathbb{Z}_{\geq 0}\)-graded ring by \(\deg x_i = w_i\).

**Definition 4.1 (Graded matrix factorization [37, 68], [53, Section 3.1]).** — A graded matrix factorization of \(W\) is a collection \((E^i, \delta_i)_{i \in \mathbb{Z}}\) of finitely generated graded free \(R\)-modules \(E^i\) and degree-zero homomorphisms \(\delta_i \in \text{Hom}_{\text{gr}-R}(E^i, E^{i+1})\)

\[
\ldots \xrightarrow{\delta_{i-1}} E^0 \xrightarrow{\delta_0} E^1 \xrightarrow{\delta_1} E^2 \xrightarrow{\delta_2} E^3 \xrightarrow{\delta_3} \ldots
\]

such that it is 2-periodic up to the shift of grading

\[
E^{i+2} = E^i(d), \quad \delta_{i+2} = \delta_i(d)
\]

and that \(\delta_{i+1} \circ \delta_i = W \cdot \text{id}_{E^i}: E^i \to E^{i+2} = E^i(d)\) for all \(i\). This is equivalent to the data \(E^0, E^1, \delta_0 \in \text{Hom}_{\text{gr}-R}(E^0, E^1), \delta_1 \in \text{Hom}_{\text{gr}-R}(E^1, E^0(d))\) such that \(\delta_1 \circ \delta_0 = W \cdot \text{id}_{E^0}\) and \(\delta_0(d) \circ \delta_1 = W \cdot \text{id}_{E^1}\). These data are denoted also by \((E, \delta_E)\), where

\[
E := E^0 \oplus E^1, \quad \delta_E := \begin{pmatrix} 0 & \delta_1 \\ \delta_0 & 0 \end{pmatrix}: E \to E \quad \text{satisfying} \quad \delta_E^2 = W \cdot \text{id}_E.
\]

These objects form a dg category as follows. Consider the graded matrix factorizations \(\overline{E} = (E^i, \delta_i)_{i \in \mathbb{Z}}\) and \(\overline{F} = (F^i, \delta'_i)_{i \in \mathbb{Z}}\); the space of homomorphisms is defined to be the \(\mathbb{Z}\)-graded vector space

\[
\text{Hom}^\bullet(\overline{E}, \overline{F}) = \{(f_n)_{n \in \mathbb{Z}} \mid f_n \in \text{Hom}_{\text{gr}-R}(E^n, F^{n'}) \cdot f_{n+2} = f_n(d)\}
\]

equipped with the differential

\[
(df)_n = \delta'_n \circ f_n - (-1)^n f_{n+1} \circ \delta_n, \quad f \in \text{Hom}^\bullet(\overline{E}, \overline{F}).
\]

The homotopy category of the above dg category is denoted by \(\text{MF}_{\mu_d}^\text{gr}(W)\). It is a triangulated category.

**Remark 4.2 ([56, Section 4.4]).** — The lower index in the notation \(\text{MF}_{\mu_d}^\text{gr}(W)\) emphasizes the fact that a graded matrix factorization is automatically \(\mu_d\)-equivariant. The \(\mu_d\)-action on \(R\) is defined by \(\zeta \cdot x_i = \zeta^{-w_i} x_i\), where \(\zeta = \exp(2\pi i / d) \in \mu_d\). For a graded matrix factorization \(\overline{E} = (E^i, \delta_i)_{i \in \mathbb{Z}}\), we define the \(\mu_d\)-action on \(E^i\) by \(\zeta \cdot e = \zeta^{-w} e\) for \(e \in (E^i)_n\). Then the \(R\)-module \(E^i\) is \(\mu_d\)-linearized and \(\delta_i\) is \(\mu_d\)-equivariant.
We introduce a graded Koszul matrix factorization (see [8, Section 2] for the ungraded case).

**Definition 4.3 (Graded Koszul matrix factorization).** — Suppose that \( W \) is of the form \( \sum_{i=1}^{N} a_i b_i \) for homogeneous elements \( a_i, b_i \in \mathbb{R} \) such that \( \deg(a_i) + \deg(b_i) = d \). Let \( V \) be the graded vector space \( \bigoplus_{i=1}^{N} \mathbb{C} e_i \) with \( \deg(e_i) = -\deg(a_i) \). For \( q \in \mathbb{Z} \), the graded Koszul matrix factorization \( \{a, b\}_q \) is defined by the data

\[
E^i = \bigoplus_{k=0, \ldots, N} \mathbb{R} \otimes \left( \bigwedge^k V \right) \left( \frac{d(i-k)}{2} + q \right),
\]

\[
\delta_i = \delta'_i + \delta''_i : E^i \to E^{i+1}, \quad i \in \mathbb{Z},
\]

where \( \delta'_i, \delta''_i \) are the Koszul differentials:

\[
\delta'_i = \sum_{j=1}^{N} a_j e_j \wedge, \quad \delta''_i = \sum_{j=1}^{N} b_j (e_j^*).\]

Observe that the grading is shifted so that the map \( \delta_i \) preserves the degree. Note also that \( \{a, b\}_q = \{a, b\}_0(q) \).

**4.1.1. Hirzebruch-Riemann-Roch Theorem.** — For graded matrix factorizations \( \overline{E}, \overline{F} \) of \( W \), we write \( \chi(\overline{E}, \overline{F}) \) for the Euler characteristic

\[
\sum_{i \in \mathbb{Z}} (-1)^i \dim H^i(\text{Hom}^*(\overline{E}, \overline{F}), d).
\]

This can be computed via Hirzebruch-Riemann-Roch (HRR) due to Walcher [68] and Polishchuk-Vaintrob [56] for \( G \)-equivariant matrix factorizations. To this effect we need to introduce the Chern character taking values in the orbifold Jacobi space \( \bigoplus_{k=0}^{d-1} \Omega(W_k)^{k^*} \), which is identified with the FJRW state space by Proposition 2.1. Let \( x_{j_1}, \ldots, x_{j_N} \) denote the co-ordinates of the \( \xi \)-fixed part \( (\mathbb{C}^N)_\xi \) where \( \xi = e^{2\pi i/d} \). For a graded matrix factorization \( \overline{E} = (E, \delta_E) \), we define [56, Theorem 3.3.3]

\[
\text{ch}(\overline{E}) := \bigoplus_{k=0}^{d-1} \left[ \text{str}_R \left( \partial_{j_1} \delta_E \circ \partial_{j_2} \delta_E \circ \cdots \circ \partial_{j_N} \delta_E \circ \xi^k \right) \right] \left( \mathbb{C}^N \right)_{\xi} dx_{j_1} \wedge \cdots \wedge dx_{j_N}.
\]
Here we take a free basis of $E = E^0 \oplus E^1$ over $R = \mathbb{C}[x_1, \ldots, x_N]$ and regard $\delta_E$ as a matrix with entries in $R$; the supertrace $\text{str}_R(f)$ of an operator $f \in \text{End}_R(E)$ is defined to be $\text{tr}(f_{0,0}) - \text{tr}(f_{1,1})$, where $f_{\sigma,\sigma} : E^\sigma \to E^\sigma$, $\sigma = 0, 1$ are the components of $f$. The right hand side are meant to be the class in $\bigoplus_{d = 0}^{N - 2} \Omega(W_k)$ and lies in the $\mu_d$-invariant part. This is independent of the choice of a co-ordinate ordering or the choice of a basis of $E$.

**Remark 4.4.** — Let $\text{ch}(\bar{E})_k$ denote the $\Omega(W_k)^{\mu_d}$ component of $\text{ch}(\bar{E})$. For a graded matrix factorization $\bar{E}$, one can see that $\text{ch}(\bar{E})_k$ vanishes if $N_k$ is odd and $\text{ch}(\bar{E})_k$ is of degree $(N_k/2)d$. In terms of the Hodge decomposition, the component $\text{ch}(\bar{E})_k$ has Hodge type $(N_k/2, N_k/2)$.

**Example 4.5.** — For a general weighted homogeneous polynomial $W$, we can write $W = \sum_{j=1}^{N} a_j b_j$ with $a_j = q_j \partial W$, $b_j = x_j (q_j := w_j/d)$. The Chern character of the graded Koszul matrix factorization $\{a, b\}$ of $W$ is supported on the narrow sector. In fact, by a direct calculation, we obtain

$$
\text{ch} \left( \{a, b\} \right) = \bigoplus_{k \in \text{Nar}} \phi_k \left( \prod_{j=1}^{N} (1 - \zeta^{-w_j k}) \right) \phi_{k-1}.
$$

See [56, Proposition 4.3.4] where $\{a, b\}_0$ is denoted by $k^a$. (The case $q \neq 0$ follows from the case with $q = 0$ since $q$ is just a shift the grading.) These Chern characters span the narrow part $H_{\text{nar}}(W, \mu_d)$.

**Theorem 4.6** (Walcher [68, Section 5], Polishchuk-Vaintrob [56, Theorem 4.2.1]). — Let $\bar{E}, \bar{F}$ be graded matrix factorizations. The Euler characteristic $\chi(\bar{E}, \bar{F})$ is given by the formula:

$$
\sum_{k=0}^{d-1} \left( \prod_{k_{w_j} \neq 0} \frac{1}{1 - \zeta^{k w_j}} \right) (-1)^{N_k(N_k-1)} \frac{1}{d} \text{Res}_{W_k}(\text{ch}(\bar{E})_k, \text{ch}(\bar{F})_{d-k}).
$$

Here $\text{ch}(\bar{E})_k$ denotes the $\Omega(W_k)^{\mu_d}$-component of $\text{ch}(\bar{E})$.

**Proof.** — Because Polishchuk-Vaintrob considered the $G$-equivariant (ungraded) matrix factorizations over the ring of formal power series, we need check that the Euler characteristic does not change under the base change from the polynomial ring to the formal power series ring for graded matrix factorizations. Set $\hat{R} = \mathbb{C}[[x_1, \ldots, x_N]]$. Let $(E, \delta_E), (F, \delta_F)$ be graded matrix factorizations. Let $(\hat{E}, \hat{\delta}_E) = (E, \delta_E) \otimes_R \hat{R}$, $(\hat{F}, \hat{\delta}_F) = (F, \delta_F) \otimes_R \hat{R}$ be $\mu_d$-equivariant matrix factorizations over $\hat{R}$ without the $\mathbb{Z}$-grading. We have the identification as $\mathbb{Z}/2$-graded complexes:

$$
\text{Hom}^\sigma((\hat{E}, \hat{\delta}_E), (\hat{F}, \hat{\delta}_F)) = \bigoplus_{j = 0}^{\sigma} \text{Hom}^j((E, \delta_E), (F, \delta_F)), \quad \sigma \in \mathbb{Z}/2,
$$
where the completed direct sum consists of arbitrary sequences of homomorphisms bounded in the negative direction. Hence the cohomology is again the completed direct sum of the cohomology $H^j(\text{Hom}^\bullet((E, \delta_E), (F, \delta_F)))$. The HRR for the left-hand side implies the finite-dimensionality and the boundedness of the cohomology of the right-hand side, and the HRR for the right-hand side as well. \[\square\]

**Remark 4.7.** Dyckerhoff [22] identified the Hochschild homology of the category of matrix factorizations over a formal power series ring with the Jacobi space of the potential. Polishchuk-Vaintrob [56] observed that the Hochschild homology can be identified with the FJRW state space in the G-equivariant case. The Chern character naturally takes values in the Hochschild homology and the Riemann-Roch formula was derived in the categorical framework in [56].

### 4.2. Orlov equivalence.

Under the Calabi-Yau condition $d = \sum_{j=1}^{N} w_j$, Orlov [53, Theorem 2.5] constructed the equivalence of triangulated categories

\begin{equation}
\Phi_l: \text{MF}_{\mu}^\text{gr}(W) \longrightarrow \text{D}^\text{b}(X_W)
\end{equation}

parametrized by $l \in \mathbb{Z}$. Consider a graded matrix factorization $E = (E_i, \delta_i)_{i \in \mathbb{Z}}$ of $W$

\[\cdots \rightarrow E^0 \xrightarrow{\delta_0} E^1 \xrightarrow{\delta_1} E^2 = E^0(d) \xrightarrow{\delta_2=\delta_0(d)} E^3 = E^1(d) \xrightarrow{\delta_3=\delta_1(d)} \cdots\]

and set

\[S = R/(W) = \mathbb{C}[x_1, \ldots, x_N]/(W).\]

By tensoring the above data $(E^i, \delta_i)_{i \in \mathbb{Z}}$ with $S$ over $R$, we obtain an acyclic (see Eisenbud [23] and Buchweitz [7]) complex

\[\cdots \xrightarrow{\delta_{-1} \otimes S} C^0 \xrightarrow{\delta_0 \otimes S} C^1 \xrightarrow{\delta_1 \otimes S} C^2 = C^0(d) \xrightarrow{\delta_2=\delta_0(d)} C^3 = C^1(d) \xrightarrow{\delta_3=\delta_1(d)} \cdots.\]

By construction we can extract from $C^\bullet$ a positively graded and left semiinfinite complex $L^\bullet \text{.}$ To this effect, after expressing each $C^i$ as a direct sum of $S$-modules of the form $S(k)$ for some $k \in \mathbb{Z}$, we mod out the $S$-modules of the form $S(-\epsilon)$ with $\epsilon \leq 0$. More precisely we may notice that $E^0$ and $E^1$ have the same dimension and can be expresses as

\[E^0 = \bigoplus_{1 \leq h \leq r} R(-j_h), \quad E^1 = \bigoplus_{r+1 \leq h \leq 2r} R(-j_h).\]

In this way we have $C^0 = \bigoplus_{1 \leq h \leq r} S(-j_h)$, and $C^1 = \bigoplus_{r+1 \leq h \leq 2r} S(-j_h)$ and

\[C^i = \bigoplus_{1 \leq h - 2(i/2) \leq r} S(d\lfloor i/2 \rfloor - j_h)\]
(note that $2r(i/2)$ equals 0 or $r$ according to the parity of $i$). Then, the definition of $L_0^\bullet$ reads

$$L_0^i = \bigoplus_{1 \leq h - 2r(i/2) \leq r \atop d[i/2] - j_h < l} S(d[i/2] - j_h).$$

Since $C^\bullet$ is acyclic, $L_0^\bullet$ is represented by a bounded complex of coherent sheaves. For simplicity, we stated the definition of the positively graded complex $L_0^\bullet$. For any $l \in \mathbb{Z}$, we can define $L_l^\bullet$ as

$$L_l^i = \bigoplus_{1 \leq h - 2r(i/2) \leq r \atop d[i/2] - j_h < l} S(d[i/2] - j_h).$$

This amounts to extracting from each $C_i$, only the $S$-modules of the form $S(-e)$ with $e > l$. We have the following statement. (We stress that the equivalence of categories holds only under the CY condition $d = \sum_{j=1}^N w_j$, which we assume throughout the paper.)

**Proposition 4.8 (Herbst-Hori-Page [34, Section 10.6, (10.56–58)]). — The Orlov equivalence**

$$\Phi_l : \text{MF}_{\mu_d}^\text{gr}(W) \to D^b(X_W)$$

for $l \in \mathbb{Z}$ assigns to $(E, \delta_E) \in \text{MF}_{\mu_d}^\text{gr}(W)$ the left semiinfinite complex (52)

$$\Phi_l(E, \delta_E) = L_l^\bullet \in D^b(X_W).$$

**Here the graded module $S(k)$ in $L_l^\bullet$ is identified with the sheaf $O(k)$ on $X_W$.**

**Remark 4.9. —** We point out that there are two presentations of $\Phi_l(E, \delta_E)$ in the derived category. Because the complex $C^\bullet$ is acyclic, the left semiinfinite complex $L_l^\bullet$ can be equivalently represented by the (complementary) right semiinfinite complex $(L_l^\bullet)^*[1]$, where

$$\begin{align*}
(L_l^i)^{[1]} = \bigoplus_{1 \leq h - 2r(i/2) \leq r \atop d[i/2] - j_h \geq l} S(d[i/2] - j_h).
\end{align*}$$

**Remark 4.10 (Herbst-Hori-Page brane transportation). —** Although we will not use this in the rest of the paper, we should mention that Orlov functors $\Phi_l$ can be constructed, for $\tilde{R} = R[\rho]$ and $\tilde{W} = \rho W$, by lifting the $\mu_d$-action to a $\mathbb{C}^\times$-action and by obtaining in this way a graded and $\mathbb{C}^\times$-equivariant matrix factorization in $\text{MF}_{\mathbb{C}^\times}^\text{gr}(\tilde{W})$. Clearly $\mu_d$-actions are not uniquely lifted to $\mathbb{C}^\times$-actions; we need an extra datum of an integer parameter $l$. This point of view due to Herbst, Hori, and Page explains the presence of several Orlov functors $\Phi_l$ for $l \in \mathbb{Z}$. From $\text{MF}_{\mathbb{C}^\times}^\text{gr}(\tilde{W})$ a natural functor leads to $D^b(X_W)$, see [34].
We apply Orlov’s functor $\Phi_l$ to the graded Koszul matrix factorization $\{a, b\}_q$ from Example 4.5 (see also Definition 4.3).

**Proposition 4.11.** — The image via $\Phi_l$ of the graded matrix factorization $\{a, b\}_q$ in Example 4.5 is represented by the complex on $\mathbb{X}_W$

$$\bigoplus_{j_1 < j_2 < \ldots < j_r, \sum_{a=1}^r w_{j_a} \leq m} \mathcal{O}\left(l + m - \sum_{a=1}^r w_{j_a}\right) e_{j_1} \wedge \cdots \wedge e_{j_r} [r + 1 + 2l]$$

equipped with the Koszul differential $\delta''_{i} = \sum_{j=1}^N \chi_i(e_i^*)$. Here $t$ and $m$ denote the integer quotient and remainder of $q - l$ divided by $d$.

**Proof.** — Write $(E^i, \delta_i)_{i \in \mathbb{Z}} = \{a, b\}_q$. Let us consider $E^i \otimes_R S$

$$\bigoplus_{r \in \mathbb{Z}, i \in (2)} S \otimes \left(\bigwedge^r V\right)\left(\frac{d(i - r)}{2} + q\right), \tag{54}$$

where $V = \bigoplus_{j=1}^N C_{e_j}$ with $\deg(e_j) = - \deg(a_j) = w_j - d$. Each summand is of the form

$$\bigoplus_{j_1 < \ldots < j_r} S \left(- \sum_{a=1}^r \deg(e_{j_a}) + \frac{d(i - r)}{2} + q\right)$$

$$= \bigoplus_{j_1 < \ldots < j_r} S \left(- \sum_{a=1}^r w_{j_a} + \frac{d(i + r)}{2} + q\right). \tag{55}$$

By Proposition 4.8 and Remark 4.9 we can regard the image via $\Phi_l$ of the Koszul matrix factorization $\{a, b\}_q$ as the (complementary) right semiinfinite complex $(L_i)[1]$. The terms $S(h)$ appearing in the above formula contribute to $L_i$ if and only if $h \geq l$; therefore we consider the inequality

$$h = - \sum_{a=1}^r w_{j_a} + \frac{i + r}{2} d + q \geq l,$$

which can be rewritten as (using $q - l = td + m$)

$$m + td + \frac{i + r}{2} d \geq \sum_{a=1}^r w_{j_a}.$$

Since $\sum_{a=1}^r w_{j_a}$ lies in $\{0, \ldots, d\}$ by the CY condition $d = \sum_{j=1}^N w_j$, we deduce that $h \geq l$ if and only if either we have $(i + r)/2 > -t$ (note $i + r$ is even by (54)) or we have $m \geq \sum_{a=1}^r w_{j_a}$ alongside with $(i + r)/2 = -t$. 
Let us consider all terms of (55) for which \((i + r)/2 > -t\). Then, the summand of (55) attached to \(j_1 < \cdots < j_r\) is of the form \(S(l + n - \sum_{a=1}^{r'} w_{ja})\) with \(n \geq d\). Such summands with fixed \(n\) form an exact sequence \(E_n^\bullet\) on \(X_W\):

\[
E_n^\bullet: \ O(l + n) \xleftarrow{\delta_k^\mu} \bigoplus_j O(l + n - w_j) \xleftarrow{\delta_k^\mu} \bigoplus_{j_1 < j_2} O(l + n - w_{j_1} - w_{j_2}) \xleftarrow{\delta_k^\mu} \cdots \xleftarrow{\delta_k^\mu} \bigoplus_j O\left(l + n - \sum_{j' \neq j} w_{j'}\right) \xleftarrow{\delta_k^\mu} O(l + n - d)
\]

where we wrote \(O(h)\) for \(S(h)\) following Proposition 4.8. Therefore, all together, the sum

\[
\bigoplus_{i \equiv r(2)} S \otimes \left(\bigwedge_{(i+r)/2 > -t}^r V\right) \left(\frac{d(i - r)}{2} + q\right)
\]

gives an acyclic subcomplex of \((L_i^\mu)^\bullet\). It is acyclic because it can be written as a successive extension by the acyclic complexes \(E_n^\bullet\) of a complex supported on arbitrarily high homological degrees. The quotient of \((L_i^\mu)^\bullet\) by this acyclic subcomplex consists of terms of (55) with \((i + r)/2 = -t\) and \(\sum_{a=1}^{r'} w_{ja} \leq m\). The conclusion follows. (Recall that we need to take the shift \((L_i^\mu)^\bullet[1]\) by 1.)

4.3. Twisted I-functions and Mellin-Barnes continuation. — We provide two parallel discussions of the twisted I-functions for GW and FJRW theories. We show that the two I-functions satisfy the same Picard-Fuchs equation under a coordinate change. We compute the connection matrix between the two I-functions (or more precisely the \(\mathcal{H}\)-functions) using the Mellin-Barnes method of analytic continuation.

On both sides we systematically work with the \(\mathbb{C}^\times\)-twisted theories. On the Landau-Ginzburg side we already discussed the \(\mathbb{C}^\times\)-twisted FJRW theory Section 3.4 over the extended state space; its non-equivariant limit, followed by projection to the narrow state space, encodes the genus zero correlator of FJRW theory. The counterpart on the Calabi-Yau side is the \(\mathbb{C}^\times\)-twisted theory of \(\mathbb{P}(w)\), twisted by \(O(d)\). It is treated and computed in genus zero in [19]; again, the non-equivariant limit, followed by the projection to the ambient part \(H_{\text{amb}}(X_W)\) of the state space yields the genus-zero correlators in GW theory. (See Section 3.5 for a review.)

4.3.1. The \(\mathbb{C}^\times\)-twisted I-functions. — Recall the \(\mathbb{C}^\times\)-twisted I-function (40) with \(N\) equivariant parameters \(\lambda_1, \ldots, \lambda_N\). Here, without loss of information from the point of view of non-equivariant theory, we can impose the conditions \(\lambda_j = -q_j \lambda\) for all \(j\) with a single equivariant parameter \(\lambda\) (as in Section 3.5.1). The \(\mathbb{C}^\times\)-twisted I-function in FJRW
theory is given by:

\[ \Gamma_{\text{FJRW}}^{\text{tw}}(u, z; \lambda) = z \sum_{k \in \mathbb{Z}_{\geq 1}} u^k \prod_{j=1}^N \prod_{0 < b < kq_j, (b) = (kj)} \frac{(-q_j \lambda - bz)}{(-bz)} \phi_{k-1}. \]

Here the index \( k - 1 \) of \( \phi_{k-1} \) is reduced modulo \( d \) within the range \( \{0, \ldots, d - 1\} \). This takes values in the extended state space \( H_{\text{ext}}(24) \).

In GW theory, the \( \epsilon_{\mathbb{C}^{\times}} \)-twisted I-function was computed in [19]. It is given by:

\[ \Gamma_{\text{GW}}^{\text{tw}}(v, z; \lambda) = z^e \log v/z \sum_{n \in \mathbb{Q}_{\geq 0}} v^n \prod_{j=1}^N \prod_{0 < b \leq dn_j, (b) = (w_j n_j) (w_j n_j - 1)} \frac{(dp + \lambda + bz)}{(-bz)} 1_{\{0\}}. \]

This encodes the \( \epsilon_{\mathbb{C}^{\times}} \)-twisted GW invariants of \( \mathbf{P}(w) \), twisted by the line bundle \( \mathcal{O}(d) \), and takes values in \( H_{\text{CR}}(\mathbf{P}(w)) \).

These twisted I-functions \( \Gamma_{\text{FJRW}}^{\text{tw}}(u, z) \) and \( \Gamma_{\text{GW}}^{\text{tw}}(v, z) \) are convergent respectively on the regions \( \{|u| < v^{-1/d}\} \) and \( \{|v| < v_c\} \), where \( v_c := d^{-d} \prod_{j=1}^N w_j \), see Lemma 5.10.

4.3.2. Picard-Fuchs equations. — The I-function \( \Gamma_{\text{FJRW}}^{\text{tw}} \) is a solution of the Picard-Fuchs equation

\[
(57) \quad u^d \prod_{j=1}^N \prod_{c=0}^{w_j-1} (w_j z D_u - q_j z - cz) - v \prod_{c=1}^d (-z D_v + cz) 1 = 0,
\]

for \( D_u = u(\partial/\partial u) \). The I-function \( \Gamma_{\text{GW}}^{\text{tw}} \) is a solution of the Picard-Fuchs equation

\[
(58) \quad \prod_{j=1}^N \prod_{c=0}^{w_j-1} (w_j z D_v - cz) - v \prod_{c=1}^d (dz D_v + \lambda + cz) 1 = 0
\]

for \( D_v = v(\partial/\partial v) \).

Under the change of variable \( u = v^{-1/d} \) and conjugation with the operator \( u^{-\lambda/z} = v^{\lambda/dz} \), the two equations coincide. This happens because we have \( dD_v = -D_u \) and \( v^{-\lambda/dz} \circ (dz D_v) \circ v^{\lambda/dz} = dz D_v + \lambda \). In particular the limits for \( \lambda \to 0 \) match under \( v = u^{-d} \). (We remedy the discrepancy of the equivariant Picard-Fuchs equations by introducing the unit co-ordinate \( \tilde{t} \) (or \( s^0 \)) later in Section 5.2.) The components of each of the I-functions give a basis of solutions to the Picard-Fuchs equation for generic \( \lambda \) (cf. Proposition 5.11, Lemma 5.15 and (69)).
4.3.3. The $\mathcal{H}$-functions. — We introduce a constant linear transform of the I-function, the $\mathcal{H}$-function, which is more compatible with the $\hat{\Gamma}$-integral structure in Section 2.4.4. The relevance of such hypergeometric series in homological mirror symmetry was observed by Horja [38], Hosono [39] and Borisov-Horja [5]. The $\mathcal{H}$-function is defined by the relation\(^15\) (cf. (22)):

\[
I^{tw}(x, z; \lambda) = z^{-\text{Gr}}\left((2\pi i)^{\text{deg}} \hat{\mathcal{H}}^{tw}(x, z; \lambda)\right).
\]

Here the operators $\hat{\Gamma}^{tw}$, $\text{Gr}$, $\text{deg}_0$ in the respective theory are defined as follows: In the twisted FJRW theory, the twisted Gamma class $\hat{\Gamma}^{tw}_{\text{FJRW}}$ operating on the extended state space $H_{\text{ext}}$ is defined to be

\[
\hat{\Gamma}^{tw}_{\text{FJRW}} := \bigoplus_{k=0}^{d-1} \prod_{i=1}^{N} \Gamma(1 - (kq_i) - q_i \xi), \quad \xi = \lambda/z.
\]

In the twisted GW theory, the twisted Gamma class $\hat{\Gamma}^{tw}_{\text{GW}}$ operating on $H_{\text{CR}}(P(w))$ is defined to be

\[
\hat{\Gamma}^{tw}_{\text{GW}} := \bigoplus_{f \in \mathcal{F}} \prod_{i=1}^{N} \Gamma(1 - (f w_i) + w_i p) \Gamma(1 + \xi + dp), \quad \xi = \lambda/z.
\]

The non-equivariant limits $\lambda \to 0$ are well-defined and induce $\hat{\Gamma}^{tw}_{\text{FJRW}}$ and $\hat{\Gamma}^{tw}_{\text{FJRW}}$ in Definition 2.17 under the projection to the original state spaces. The grading operator $\text{Gr}$ on $H_{\text{ext}}$ or on $H_{\text{CR}}(X_W)$ is given by

\[
\text{Gr}(T_i) = \frac{\text{deg} T_i}{2} T_i
\]

where “deg” denotes the degree defined in (25) for FJRW theory and the age-shifted degree of orbifold cohomology classes of $P(w)$ for GW theory. The “bare” degree operator $\text{deg}_0$ on $H_{\text{ext}}$ or on $H_{\text{CR}}(P(w))$ is defined by (cf. Definition 2.19)

\[
\text{deg}_0(\phi_k) = -2\phi_k \quad \text{for twisted FJRW theory;}
\]

\[
\text{deg}_0(\rho^n 1_j) = 2n(\rho^n 1_j) \quad \text{for twisted GW theory.}
\]

On the Landau-Ginzburg side, we have

\(^15\) See Section 5.3, (82) for a precise relationship between the $\mathcal{F}$-function and the $\hat{\Gamma}$-integral structure.
\[
\Gamma_{\text{FJRW}}^{\text{tw}}(u, z; \lambda) = z^{-\text{Gr}_z} \sum_{k \in \mathbb{Z}_{\geq 1}} u^k \frac{(-1)^{k-1}}{\Gamma(k)} \prod_{j=1}^{N} \frac{\Gamma((-q_j k) - q_j \xi)}{\Gamma(1 - q_j k + \xi)} \phi_{k-1}
\]

\[
= z^{-\text{Gr}_z} \sum_{k \in \mathbb{Z}_{\geq 1}} u^k \frac{(-1)^{k-1}}{\Gamma(k)} \frac{1}{\prod_{j:k \in \mathbb{Z}} (-q_j \xi)} \times \prod_{j=1}^{N} \frac{\Gamma(1 - (q_j k) - q_j \xi)}{\Gamma(1 - q_j (k + \xi))} \phi_{k-1}
\]

\[
= z^{-\text{Gr}_z} \hat{\Gamma}_{\text{FJRW}}^{\text{tw}} \left( (2\pi \mathbf{i})^{\text{deg}_0} \right) \hat{\mathcal{S}}_{\text{FJRW}}^{\text{tw}}(u, z; \lambda),
\]

where \( \xi := \lambda/z \) and

\[
\mathcal{S}_{\text{FJRW}}^{\text{tw}}(u, z; \lambda) = z \sum_{k \in \mathbb{Z}_{\geq 1}} \frac{(-1)^{k-1}(2\pi \mathbf{i})}{\Gamma(k)} \prod_{j:k \in \mathbb{Z}} (-q_j \xi) \prod_{j=1}^{N} \frac{\Gamma(1 - q_j (k + \xi))}{\Gamma(1 - q_j k + \xi)} \phi_{k-1}.
\]

On the Calabi-Yau side, we have

\[
\Gamma_{\text{GW}}^{\text{tw}}(v, z; \lambda) = z^{-\text{Gr}_z} e^{\delta \log v} \sum_{n \in \mathbb{Q}_{\geq 0} \atop j, n \in \mathbb{Z}} v^n \frac{\Gamma(1 + dp + \xi + dn)}{\Gamma(1 + dp + \xi)}
\]

\[
\times \prod_{j=1}^{N} \frac{\Gamma(1 + w_j p - (-w_j n))}{\Gamma(1 + w_j p + w_j n)} (\mathbf{1}_{(-n)}
\]

\[
= z^{-\text{Gr}_z} \hat{\Gamma}_{\text{GW}}^{\text{tw}} \left( (2\pi \mathbf{i})^{\text{deg}_0} \right) \hat{\mathcal{S}}_{\text{GW}}^{\text{tw}}(v, z; \lambda),
\]

where \( \xi := \lambda/z \) and

\[
\mathcal{S}_{\text{GW}}^{\text{tw}}(v, z; \lambda) = z^{\frac{\delta \log v}{2\pi \mathbf{i}}} \sum_{n \in \mathbb{Q}_{\geq 0} \atop j, n \in \mathbb{Z}} v^n \frac{\Gamma(1 + d \frac{p}{2\pi \mathbf{i}} + \xi + dn)}{\prod_{j=1}^{N} \Gamma(1 + w_j \frac{p}{2\pi \mathbf{i}} + w_j n)} \mathbf{1}_{(-n)}.
\]

**4.3.4. Mellin-Barnes analytic continuation.** — The function \( \mathcal{S}_{\text{GW}}^{\text{tw}}(v, z; \lambda) \) is convergent and analytic on the region \( \Re(\log v) < \log v_c \), where \( v_c := d^{-d} \prod_{j=1}^{N} w_j^{-w_j} \) is the singularity of the Picard-Fuchs equation \( (58) \). Similarly, \( \mathcal{S}_{\text{FJRW}}^{\text{tw}}(u, z; \lambda) \) is convergent and analytic on the region \( \Re(\log u) < -(\log v_c)/d \). Let \( \hat{\mathcal{M}}^\circ \) denote the \( (\log v) \)-plane minus the singularities of the Picard-Fuchs equation:

\[
\hat{\mathcal{M}}^\circ = \mathcal{C}_{\log v} \setminus \{ \log v_c + 2l\pi \mathbf{i} \mid l \in \mathbb{Z} \}.
\]

Under the identification \( \log v = -d \log u \), we regard \( \mathcal{S}_{\text{GW}}^{\text{tw}} \) as a single-valued function in the left-half of \( \hat{\mathcal{M}}^\circ \) and \( \mathcal{S}_{\text{FJRW}}^{\text{tw}} \) as a single-valued function on the right-half of \( \hat{\mathcal{M}}^\circ \). Let
\( \gamma_l \subset \tilde{\mathcal{M}} \) be a path from the large radius limit (\( \Im(\log v) = 0, \Re(\log v) \ll 0 \)) to the LG limit (\( \Im(\log v) = 0, \Re(\log v) \gg 0 \)) which passes through the “window” [\( \log v_c + 2(l - 1)\pi \hat{b}, \log v_c + 2(l + 1)\pi \hat{b} \)]. See Figure 3. We consider analytic continuation along the path \( \gamma_l \).

We rewrite \( H_{\text{tw}}^{\text{GW}} \) by expressing the running index \( n \) as an element of \( \tilde{\mathcal{H}} + \mathbb{Z}_{\geq 0} \). For \( f \in \tilde{\mathcal{H}} \), we adopt the notation \( f = \langle 1 - f \rangle \). We get

\[
H_{\text{tw}}^{\text{GW}}(v, z; \lambda) = z \sum_{f \in \tilde{\mathcal{H}}} \sum_{k \in \mathbb{Z}_{\geq 0}} \frac{\Gamma(1 + \xi + d \frac{p}{2\pi i} + d\overline{f} + dk)}{\prod_{j=1}^{N} \Gamma(1 + w_j \frac{p}{2\pi i} + w_j\overline{f} + w_jk)} \cdot \frac{1}{v^{\frac{p}{2\pi i} + \overline{f} + s}}.
\]

During the analytic continuation, we regard \( p \) as a small complex number and think of the \( \tilde{\mathcal{H}} \)-function as a scalar valued function. At the end of the calculation, we take the Taylor expansion in \( p \) and replace \( p \) with the hyperplane class. In this way we get analytic continuation of a cohomology-valued function. We write the sum over \( \mathbb{Z}_{\geq 0} \) as a sum of residues:

\[
z \sum_{f \in \tilde{\mathcal{H}}} \sum_{k \in \mathbb{Z}_{\geq 0}} \text{Res}_{s=\lambda} ds \quad \times \left( \frac{\Gamma(s) \Gamma(1 - s)}{\prod_{j=1}^{N} \Gamma(1 + w_j \frac{p}{2\pi i} + w_j\overline{f} + w_jk)} e^{-(2l-1)\pi \hat{b} s} \right).
\]

Here \( l \in \mathbb{Z} \) is the index of the path \( \gamma_l \). Consider the contour integrals along the path of Figure 4 of each 1-form in the above expression. The integrals are absolutely convergent (and define analytic functions of \( v \)) if \( |\Im(\log v) - (2l - 1)\pi| < \pi \) (see e.g. [38, Lemma 3.3]). This condition is satisfied when \( \log v \) is along (the middle part of) the path \( \gamma_l \). When \( |v| < v_c \), we can close the contour to the right and obtain the above sum of residues. On the other hand, if \( |v| > v_c \), we can close the contour to the left and obtain the sum of residues at \( s = -m (m \in \mathbb{Z}_{\geq 1}) \) plus the sum of residues at

\[
s = -(\frac{1 + \xi}{d} + \frac{p}{2\pi i} + \overline{f}) \quad \text{for } k \in \mathbb{Z}_{\geq 0}.
\]
The sum of these residues gives

\begin{equation}
- z \sum_{f \in \mathfrak{D}} \sum_{m=1}^{\infty} \frac{\Gamma(1 + \xi + d(\frac{p}{2\pi \mathbb{I}} + \bar{f} - m))}{\prod_{j=1}^{N} \Gamma(1 + w_j(\frac{p}{2\pi \mathbb{I}} + \bar{f} - m))} e^{(\frac{p}{2\pi \mathbb{I}} + \bar{f} - m) \log v} \prod_{j=1}^{N} \frac{1}{\Gamma(1 + w_j)} \left( 1 - q_j(1 + \xi + k) \right)
\end{equation}

Here the overall minus sign appears because the contour closed to the left encloses each pole clockwise. We also used the co-ordinate change \( \log u = -\log v/d \).

We now regard \( p \) as the hyperplane class on \( \mathbf{P}(w) \). The first term of (63) vanishes in cohomology because the class

\[ \prod_{j: w_j f \in \mathbb{Z}} \frac{1}{\Gamma(1 + w_j(\frac{p}{2\pi \mathbb{I}} + \bar{f} - m))} = \mathcal{O}(p) \]

is zero on the sector \( \mathbf{P}(w)_f \). (Note that \( \mathbf{P}(w)_f \) is of dimension \( \# \{ j \mid w_j f \in \mathbb{Z} \} - 1 \).) By shifting the index \( k \) by 1 and using \( \sin(x) = (e^{ix} - e^{-ix})/2i \), we can rewrite the second term of (63) as

\[ z \sum_{f \in \mathfrak{D}} \sum_{k=0}^{\infty} \frac{1}{d} \frac{(\xi + d(\frac{p}{2\pi \mathbb{I}} + \bar{f} + \frac{k}{2}))^l}{(k - 1)! \prod_{j=1}^{N} \Gamma(1 - q_j(k + \xi))} \cdot e^{(\frac{p}{2\pi \mathbb{I}} + \bar{f} + \frac{k}{2})} \cdot (-1)^{k-1} u^{\xi + k} \]

This expression is regular at \( p = 0 \) and can be regarded as an \( H_{CR}(\mathbf{P}(w)) \)-valued function. This is the analytic continuation of \( S_{tGW}^{\mathbb{Z}} \) along the path \( \gamma_t \). Comparing this with \( S_{tFW}^{\mathbb{Z}}(60) \), we have the following proposition:
Proposition 4.12. — Define a linear transformation \(U_{tw}^{l}: H_{ext} \to H_{CR}(P(w))\) depending on \(l \in \mathbb{Z}\) and the parameter \(\xi = \lambda/z\) by

\[
U_{tw}^{l}(\phi_{k-1}) = \frac{1}{d} \sum_{f \in \mathfrak{F}} \frac{1}{1 - \sum_{j:q_j \in \mathbb{Z}} (-q_j \xi)} \left( \xi^{i} e^{2\pi i (\gamma + \xi)} \right)^l \prod_{j:q_j \in \mathbb{Z}} (-q_j \xi), \quad k = 1, \ldots, d.
\]

Then we have

\[
u^{-\xi} \left( \mathfrak{F}_{GW}^{tw} \right)_{\text{continued}} = U_{tw}^{l} \left( \mathfrak{F}_{FJR W}^{tw} \right),
\]

where \(\mathfrak{F}_{GW}^{tw} \) is the analytic continuation of \(\mathfrak{F}_{GW}^{tw}\) along the path \(\gamma_l\).

Remark 4.13. — By Proposition 4.12 and (59), we can find the connection matrix of the twisted I-functions. We have \(u^{-\xi}(\mathfrak{F}_{GW}^{tw})_{\text{continued}} = \tilde{U}_{tw}^{l}(\mathfrak{F}_{FJR W}^{tw})\) for the transformation

\[
\tilde{U}_{tw}^{l} = z^{-Gr} \circ \hat{\Gamma}_{GW} \circ (2\pi i)^{-\frac{d \dim_P(w)}{2\pi i}} \circ U_{tw}^{l} \circ (2\pi i)^{-\frac{d \dim_P(w)}{2\pi i}} \circ \hat{\Gamma}_{FJR W} \circ (2\pi i)^{-\frac{d \dim_P(w)}{2\pi i}} \circ z^{Gr}.
\]

The non-equivariant limit of this induces a linear transformation between the Givental symplectic vector spaces of FJRW theory and GW theory. This is the symplectic transformation computed in [13] for a quintic.

4.4. The non-equivariant limit and Orlov equivalence. — Here we show that the non-equivariant limit of \(U_{tw}^{l}\) exists and descends to a linear transformation between the narrow and the ambient part state spaces. We show that it matches with the numerical Orlov equivalence.

4.4.1. The narrow-to-ambient linear transformation.

Proposition 4.14. — The non-equivariant limit \(\lambda \to 0\) of \(U_{tw}^{l}\) exists. We have

\[
\lim_{\lambda \to 0} \left( U_{tw}^{l}(\phi_{k-1}) \right) = \begin{cases} \frac{\gamma^i}{d} \sum_{f \in \mathfrak{F}} \left( \xi^{i} e^{2\pi i (\gamma + \xi)} \right)^l \prod_{j:q_j \in \mathbb{Z}} (-q_j \xi) & \text{for } k \in \text{Nar}; \\ -p^{N_k-1} \frac{1}{d} \sum_{f \in \mathfrak{F}} \xi^{i} e^{2\pi i (\gamma + \xi)} \prod_{j:q_j \in \mathbb{Z}} (-q_j \xi) & \text{for } k \notin \text{Nar}, \end{cases}
\]

where \(k = 1, \ldots, d\) and \(N_k := 1 + \dim_P(w)_{(k/d)} = \sharp \{j \mid kq_j \in \mathbb{Z} \} \).
If \( k \notin \text{Nar}, \) (64) is not regular at \( (p, \xi) = (0, 0) \). The only non-regular term in (64) is the one with \( f = \langle k/d \rangle \) (in this case \( \zeta_k e^{2\pi i j} = 1 \)). We compute the Taylor expansion in \( p \) of such term. By an elementary computation, we have

\[
\frac{1}{e^{\phi + \frac{2\pi i f}{d}} - 1} = \sum_{n=0}^{\infty} \beta_n(\xi) p^n, \quad \beta_n(\xi) = (-1)^n \left( \frac{2\pi i \xi}{d} \right)^{-n-1} + O(\xi^{-n}).
\]

When evaluated in the cohomology group \( H(P(w)_{\mu d}) \), this Taylor series is truncated at \( n = \dim P(w)_{\mu d} = N_k - 1 \) (where we used \( f = \langle k/d \rangle \)). Therefore the factor \( \prod_{j: kq_j \in \mathbb{Z}} (\zeta_k e^{2\pi i f} - 1) \) cancels all the negative powers of \( \xi \) in \( \beta_n \). Hence \( U^w_{\mu d}(\phi_{k-1}) \) is regular at \( \xi = 0 \) and the conclusion follows. \( \square \)

We have natural projections \( H_{\text{ext}} \to H_{\text{nar}}(W, \mu_d), H_{\text{CR}}(P(w)) \to H_{\text{amb}}(X_W) \) from the state spaces of the twisted theory to the narrow/ambient part of the state spaces. We denote this projection by \( \text{pr} \). By Proposition 4.14, \( \lim_{\lambda \to 0} U^w \) descends to these projections.

**Corollary 4.15.** — Define a linear transformation \( U_f: H_{\text{nar}}(W, \mu_d) \to H_{\text{amb}}(X_W) \) by

\[
U_f(\phi_{k-1}) = \frac{1}{d} \sum_{f \in \mathbb{N}} \frac{(\zeta_k e^{\phi + 2\pi i f})^f}{\zeta_k e^{\phi + 2\pi i f} - 1} 1_f.
\]

Then we have the commutative diagram:

\[
\begin{array}{ccc}
H_{\text{ext}} & \xrightarrow{\lim_{\lambda \to 0} U^w_{\mu d}} & H_{\text{CR}}(P(w)) \\
\downarrow \text{pr} & & \downarrow \text{pr} \\
H_{\text{nar}}(W, \mu_d) & \xrightarrow{U_f} & H_{\text{amb}}(X_W)
\end{array}
\]

The operator \( U_f \) gives a connection between the non-equivariant limit of \( \mathcal{F} \)-functions, i.e. \( \mathcal{F}_{\text{GW}} = U_f(\mathcal{F}_{\text{FRW}}) \) for \( \mathcal{F}_{\phi} := \text{pr}(\lim_{\lambda \to 0} \mathcal{F}^w_{\mu d}) \).

**4.4.2.** The analytic continuation matches Orlov equivalences. — Via the Chern character, the linear transformations \( U_f \) match the Orlov equivalences \( \Phi_f \). To show this we use the explicit expression for Orlov’s equivalence for Koszul matrix factorizations (Proposition 4.11) and the equation (49) for the Chern character.

**Lemma 4.16.** — We have

\[
\frac{1}{d} \sum_{k=0}^{d-1} \zeta_k^{kj} \zeta_k^{ky} - 1 = \frac{y^{d(-j/d)}}{y^d - 1},
\]

where \( d(-j/d) \) is simply \( -j \) reduced modulo \( d \) within \( \{0, 1, \ldots, d - 1\} \).
Proof. — Note that \( \frac{1}{d} \sum_{k=0}^{d-1} \zeta^{kj} \) equals 1 if \( q \in d\mathbb{Z} \) and 0 otherwise. Thus we have

\[
\frac{1}{d} \sum_{k=0}^{d-1} \zeta^{kj} y - 1 = \frac{1}{d} \sum_{k=0}^{d-1} \sum_{n=0}^{\infty} (\zeta^k)^{j+n} y^n = \sum_{n \geq (kj+n) \in d\mathbb{Z}} y^n.
\]

The lemma follows. \( \square \)

Theorem 4.17. — Let \( U_l : H_{\text{nar}}(W, \mu_d) \to H_{\text{amb}}(X_W) \) denote the map in Corollary 4.15. For a graded matrix factorization \( E \in \text{MF}^{gr}_{\mu_d}(W) \) such that \( \text{ch}(E) \in H_{\text{nar}}(W, \mu_d) \), we have

\[ U_l(\text{inv}^* \text{ch}(E)) = \text{inv}^* \text{ch}(\Phi_l(E)). \]

Proof. — Because Chern characters of the form \( \text{ch}(\{a, b\}_q) \) in Example 4.5 span the narrow part, it suffices to show that

\[ U_l(\text{inv}^* \text{ch}(\{a, b\}_q)) = \text{inv}^* \text{ch}(\Phi_l(\{a, b\}_q)) \]

for \( q \in \mathbb{Z} \) and \( a, b \) in Example 4.5. Using (65) and (49), we get

\[
U_l(\text{inv}^* \text{ch}(\{a, b\}_q)) = U_l\left( \sum_{k \in \text{Nar}} \zeta^{-qk} (1 - \zeta^{wj}) \cdots (1 - \zeta^{wn}) \phi_{k-1} \right)
\]

\[
= \frac{1}{d} \sum_{f \in \mathfrak{S}} \sum_{k=0}^{d-1} \zeta^{-qk} (1 - \zeta^{wj}) \cdots (1 - \zeta^{wn}) \frac{\zeta^k e^{2\pi i f} - 1}{\zeta^k e^{2\pi i f} - 1} \left( \zeta^k e^{2\pi i f} \right)^{1_f}
\]

\[
= \sum_{f \in \mathfrak{S}} 1_f \sum_{j_1 < \cdots < j_r} (\zeta^l)^{-q+l+w_{j_1}+\cdots+w_{j_r}} (\zeta^l)^{j_1} (\zeta^l)^{j_2} \cdots (\zeta^l)^{j_r} (\zeta^l)^{-q-l+w_{j_1}+\cdots+w_{j_r}} ,
\]

where we set \( y := e^{2\pi i f} \). Using Lemma 4.16, we can write the coefficient of \( 1_f \) as

(66) \[
\frac{y^l}{1-y^d} \sum_{j_1 < \cdots < j_r} (-1)^{r+1} y^{d \left( \sum_{i=1}^{r} w_{j_i} - \frac{1}{d} \sum_{i=1}^{r} w_{j_i} \right)}.
\]

Let \( m \) be the remainder of \( q - l \) divided by \( d \). The sum (66) can be decomposed as

\[
\frac{y^l}{1-y^d} \left( \sum_{\sum_{i=1}^{r} w_{j_i} \leq m} (-1)^{r+1} y^{\sum_{i=1}^{r} w_{j_i} - \frac{1}{d} \sum_{i=1}^{r} w_{j_i}} + \sum_{\sum_{i=1}^{r} w_{j_i} > m} (-1)^{r+1} y^{m - \sum_{i=1}^{r} w_{j_i} + \frac{1}{d}} \right).
\]
This can be further rewritten as

\[(67) \quad \frac{y^d}{1 - y^d} \left( (1 - y^d) \sum_{j_1 < \cdots < j_r} (-1)^{r+1} y^{m - \sum_{i=1}^r w_{j_i}} + \sum_{j_1 < \cdots < j_r} (-1)^{r+1} y^{m - \sum_{i=1}^r w_{j_i} + d} \right). \]

The second summand equals

\[-\frac{y^{d+l+m}(1 - y^{w_1}) \cdots (1 - y^{w_N})}{1 - y^d}.\]

This is divisible by \( p^{2|w_j/\mathbb{Z}| - 1} \) and vanishes in \( H(P(w)/\cap X_W) \) for the dimensional reason (note that \( \dim(P(w)/\cap X_W) = 2|\{j | w_j \in \mathbb{Z}\}| - 2 \)). Finally, the first summand of \((67)\) equals the coefficient of \( 1_f \) of \( \text{inv}^* \text{ch}(\Phi_1(\{g, h\})) \) by Proposition 4.11.

\[
\square
\]

5. Construction of global D-module

This section is devoted to the proof of the main theorems in Sections 1.1 and 2.5. We construct a global D-module over the base \( \mathcal{M} = P(1, d) \setminus \{2 \text{ points}\} \) as an explicit GKZ-type differential system and show that the D-module is isomorphic to the quantum D-module of GW theory near \( v = 0 \) and to the quantum D-module of FJRW theory near \( v = \infty \). We use the mirror theorem in Section 3 and that of Coates-Corti-Lee-Tseng \[19\] (and its refinement in \[43\]).

5.1. Multi-GKZ system. — Let \( v \mapsto [1, v] \) denote the inhomogeneous co-ordinate on \( P(1, d) \) where \( v = \infty \) is the \( \mathfrak{g}_d \)-orbifold point (LG point). Using the co-ordinate \( v \), we set

\[ \mathcal{M} := P(1, d) \setminus \{0, v_c\}, \quad \mathcal{M}^c := P(1, d) \setminus \{0, v_c, \infty\} \]

where \( v_c := d^{-d} \prod_{j=1}^N w_j^{w_j} \) is the conifold point. Let \( u := v^{-1/d} \) denote the uniformizing co-ordinate centered at the LG point. In this section we introduce a GKZ-type (Gelfand-Kapranov-Zelevinskii \[27\]) hypergeometric D-module over the base \( \mathcal{M}^c \). The D-module here involves the parameter \( z \) which appears in the quantum D-module (see Section 2.4.2), and the equivariant parameter \( \lambda \) which appears in the twisted theory (see Section 3). Therefore it is defined as a sheaf over \( \mathcal{M}^c \times C_z \times C_\lambda \). Let \( \mathcal{R}^{\text{tw}} \) denote the sheaf of algebras over \( \mathcal{M}^c \times C_z \times C_\lambda \) given by the non-commutative ring of differential operators

\[ C(z, \lambda, v^\pm, (v - v_c)^{-1}, zD_v) \]

where \( D_v = v(\partial/\partial v) \). We also set

\[ B := \{(v_0, \ldots, v_N) \in \mathbb{Z}^{N+1} \mid v_i + q_i v_0 \geq 0, i = 1, \ldots, N\}, \quad q_i = w_i/d. \]
Definition 5.1. — The sheaf $\mathcal{F}^{tw}$ over $\mathcal{M}^0 \times \mathbf{G}_z \times \mathbf{G}_\lambda$ is defined to be the $\mathcal{R}^{tw}$-module generated by the symbols $\Delta_\nu$ with $\nu \in B$ subject to the relations:

\begin{align}
(dzD_\nu + \lambda + (v_0 + 1)z)\Delta_\nu &= \Delta_{\nu+\alpha_0}, \\
(w_i zD_\nu - v_i z)\Delta_\nu &= \Delta_{\nu+\epsilon_i}, & i \in \{1, \ldots, N\}, \\
v \cdot \Delta_\nu &= \Delta_{\nu+(-d, w_1, \ldots, w_N)}.
\end{align}

(68)

Here $\nu \in B$ and $\epsilon_i = (0, \ldots, 0, i, 0, \ldots, 0)$, $0 \leq i \leq N$. This defines a GKZ-type hypergeometric differential system. In fact, it is easy to see that each generator $\Delta_\nu$ satisfies the relation

\begin{align}
\left[ v \prod_{k=1}^d (dzD_\nu + \lambda + (v_0 + k)z) - \prod_{i=1}^N \prod_{k=0}^{w_i-1} (w_i zD_\nu - (v_i + k)z) \right] \Delta_\nu = 0.
\end{align}

(69)

Remark 5.2. — A multi-generated hypergeometric system similar to $\mathcal{F}^{tw}$ above appeared in the recent work of Borisov-Horja [6] (also will appear in Coates-Corti-Iritani-Tseng [18]). The $\mathcal{R}^{tw}$-submodule $\mathcal{R}^{tw}_{\Delta_0}$ of $\mathcal{F}^{tw}$ generated by $\Delta_0$ coincides with $\mathcal{F}^{tw}$ at the generic point (Lemma 5.15), but not everywhere (for instance along $z = \lambda = 0$). A closely related multi-generation phenomena of quantum cohomology was observed by Guest-Sakai [33] for a Fano hypersurface in $\mathbf{P}(w)$. It was shown in [43] that the quantum D-module of a toric Calabi-Yau hypersurface can be described by a multi-GKZ system.

Remark 5.3. — Givental’s mirror [28] (adapted to a Calabi-Yau hypersurface $X_w$ in the weighted projective space $\mathbf{P}(w)$) gives a solution to the above differential system. Let $x_0, \ldots, x_N$ be mirror $\mathbf{C}^\times$-variables subject to the relation

\begin{align}
x_0^{-d} x_1^{w_1} \cdots x_N^{w_N} = v.
\end{align}

The mirror potential $W_\lambda$ is defined by

\begin{align}
W_\lambda = x_1 + \cdots + x_N - x_0 + \lambda \log x_0.
\end{align}

Then the integrals

\begin{align}
\mathcal{I}_\nu(v) = \int x_0^{v_0} x_1^{v_1} \cdots x_N^{v_N} e^{w_\lambda/\nu} \frac{dx_0 \wedge d\log x_1 \wedge \cdots \wedge d\log x_N}{d\log v}, & \quad \nu \in B
\end{align}

satisfy the same differential relations as $\Delta_\nu$’s do. The integration cycle is contained in the torus $\{ (x_0, \ldots, x_N) \in (\mathbf{C}^\times)^{N+1} \mid x_0^{-d} x_1 \cdots x_N = v \}$ and possibly noncompact, but here we do not try to justify the integral itself. The differential relations among $\mathcal{I}_\nu(v)$ follow from a formal computation of integration by parts.
**Lemma 5.4.** — Set \( v(l) := (l, -\lfloor q_1 l \rfloor, \ldots, -\lfloor q_N l \rfloor) \in B. \) The sheaf \( F^{tw} \) is generated by \( \Delta_{v(l)}, \ l = 0, \ldots, d - 1 \) as an \( R^{tw} \)-module.

**Proof.** — For \( v = (v_0, \ldots, v_N) \in B, \) set \( l = d(v_0/d). \) Observe that

\[
v = v(l) + \sum_{i=1}^{N} (v_i + \lfloor q_i v_0 \rfloor) e_i - \left[ \frac{v_0}{d} \right] (-d, w_1, \ldots, w_N), \quad v_i + \lfloor q_i v_0 \rfloor \geq 0.
\]

The conclusion follows from this and the defining relations (68) of \( F^{tw}. \) □

The sheaf \( F^{tw} \) is a \( \mathbb{Z}_{\geq 0} \)-graded \( R^{tw} \)-module with respect to the grading

\[
\deg v = 0, \quad \deg z = \deg \lambda = \deg(zD_v) = 2,
\]

\[
\deg \Delta_v = 2(v_0 + \cdots + v_N).
\]

(Strictly speaking, the module of global sections of \( F^{tw} \) is graded, but we abuse the language since we are working over the affine base.)

**Lemma 5.5.** — Set \( \delta(l) := \frac{1}{2} \deg \Delta_{v(l)}. \) We have

(i) \( \delta(l+1) \leq \delta(l) + 1, \) \( \delta(l+d) = \delta(l). \)

(ii) \( 0 \leq \delta(l) \leq N - 1. \) We have \( \delta(l) = N - 1 \) if and only if \( l \equiv -1 \mod d. \)

**Proof.** — We have \( \delta(l) = l - \sum_{i=1}^{N} \lfloor q_i l \rfloor. \) Part (i) follows from this formula. Part (ii) follows from \( \delta(l) = \sum_{i=1}^{N} \lfloor q_i l \rfloor \leq \sum_{i=1}^{N} (1 - q_i) = N - 1. \) The equality holds iff \( l \equiv -1 \mod d/w_i \) for all \( i, \) i.e. \( l \equiv -1 \mod d. \) □

**Lemma 5.6.** — The following relations hold in \( F^{tw}: \)

(i) For \( 0 \leq l < d - 1, \) \( m = \min \{ l \leq l' \leq d - 1 \mid \delta(l') = \delta(l) + 1 \} \) exists and we have

\[
zD_v \cdot \Delta_{v(l)} \in d^{l-m} \left( \prod_{i=1}^{N} w_i^{\lfloor q_i m \rfloor - \lfloor q_i l \rfloor} \right) \Delta_{v(m)} + (z, \lambda) F^{tw}.
\]

(ii) \( zD_v \cdot \Delta_{v(d-1)} \in (z, \lambda) F^{tw}. \)

**Proof.** — The existence of \( m \) follows from Lemma 5.5. We have by (68)

\[
(dzD_v + \lambda + (l+1)z) \Delta_{v(l)} = \prod_{i=1}^{N} \prod_{\lfloor q_i l \rfloor < k \leq \lfloor q_i (l+1) \rfloor} (w_i z D_v + kz) \Delta_{v(l+1)}.
\]

Hence

\[
zD_v \cdot \Delta_{v(l)} \in d^{-1} \left( \prod_{i=1}^{N} w_i^{\lfloor q_i (l+1) \rfloor - \lfloor q_i l \rfloor} \right) (zD_v)^{\delta(l) - \delta(l+1) + 1} \Delta_{v(l+1)} + (z, \lambda) F^{tw}.
\]
If $\delta(l+1) \leq \delta(l)$, we can apply this formula recursively for $zD_v \cdot \Delta_{v(l+1)}$ in the right-hand side. In general, if $\delta(l') \leq \delta(l)$ for all $l'$ with $l < l' < m'$, we have

$$zD_v \cdot \Delta_{v(l)} \in d^{-(m'-l)} \left( \prod_{i=1}^{N} w_i^{lq_{m'}-[q,l]} \right) (zD_v)^{\delta(l)-\delta(m')+1} \Delta_{v(m')} + (z, \lambda) F^{tw}.$$ 

Taking $m'$ to be $m$, we have (i). When $l = d - 1$, we can take $m'$ to be $l + d = 2d - 1$. Then we have

$$zD_v \cdot \Delta_{v(d-1)} - v_c zD_v \cdot \Delta_{v(2d-1)} \in (z, \lambda) F^{tw}. $$

Part (ii) follows because $\Delta_{v(2d-1)} = v^{-1} \Delta_{v(d-1)}$ and $(1 - v_i/v)$ is invertible. □

**Theorem 5.7.** — The sheaf $F^{tw}$ is a free $\mathcal{O}_{M^0 \times C_x \times C_\lambda}$-module of rank $d$ with the basis $\Delta_{v(l)}$, $l = 0, \ldots, d - 1$.

**Proof.** — Let $F^{tw'}$ be the $\mathcal{O}_{M^0 \times C_x \times C_\lambda}$-submodule of $F^{tw}$ generated by $\Delta_{v(l)}$, $l = 0, \ldots, d - 1$. First we see that $F^{tw'} = F^{tw}$. We proceed by induction on the degree. The degree zero part $(F^{tw})_0$ is generated by $\Delta_0 = \Delta_{v(0)}$. Hence $(F^{tw})_0 \subset F^{tw'}$. Assume by induction that $(F^{tw})_{\leq 2k} \subset F^{tw'}$ for some $k \geq 0$. We shall show $(F^{tw})_{\leq 2(k+1)} \subset F^{tw'}$. By Lemma 5.4, it suffices to show that $zD_v \cdot \Delta_{v(l)} \in F^{tw'}$ for $0 \leq l \leq d - 1$ with $\delta(l) = k$. This follows from Lemma 5.6 and the induction hypothesis. Therefore $F^{tw} = F^{tw'}$.

As we will see in Proposition 5.11 below, $F^{tw}$ has $d$ independent solutions. This shows that the generic rank (the rank at the generic point) of $F^{tw}$ equals $d$. By the previous paragraph, $F^{tw}$ is generated by $\Delta_{v(l)}$, $l = 0, \ldots, d - 1$. Suppose we have a relation

$$\sum_{i=0}^{d-1} f_i(v, \lambda, z) \Delta_{v(l)} = 0 \text{ with } f_i \in \mathcal{O}_{M^0 \times C_x \times C_\lambda}. $$

Then $f_i$ should vanish at the generic point. Therefore $f_i = 0$. The conclusion follows. □

**5.2. Refined mirror theorem.** — We construct a basis of hypergeometric solutions of the GKZ system $F^{tw}$. Then we relate it to the fundamental solution $L^{tw}$ of the $e_{C_x \times C_\lambda}$-twisted quantum connection (see Section 3.5) in Theorem 5.12. This shows the analytic continuation of twisted quantum connections. (In this section “twisted” always means “$e_{C_x \times C_\lambda}$-twisted”.)

First we will “thicken” $F^{tw}$ by adding a new co-ordinate $l^0$. Let $\tilde{M}^0 \to M^0$ be the minimal abelian cover of $M^0$ such that $\log v$ is single-valued (see (62)). We set

$$\tilde{M} = C_{l^0} \times \tilde{M}^0,$$

where $C_{l^0}$ denotes the complex plane with co-ordinate $l^0$. Define another co-ordinate $s^0 : \tilde{M} \times C_\lambda \to C$ by

$$s^0 = l^0 - \frac{1}{d} \lambda \log v.$$
We shall use \((t^0, v; \lambda)\) and \((s^0, u; \lambda)\) as two co-ordinate systems on \(\hat{\mathcal{M}} \times \mathbf{C}_\lambda; \) \((t^0, v; \lambda)\) is a chart for GW theory and \((s^0, u; \lambda)\) is for FJRW theory. The co-ordinates \(t^0\) and \(s^0\) correspond to the identity direction of the state space via the mirror map we consider below. Let \(\hat{\mathcal{R}}^{\text{tw}}\) be the following sheaf of algebras over \(\hat{\mathcal{M}} \times \mathbf{C}_\zeta \times \mathbf{C}_\lambda;\)

\[
\hat{\mathcal{R}}^{\text{tw}} = \mathcal{O}_{\hat{\mathcal{M}} \times \mathbf{C}_\zeta \times \mathbf{C}_\lambda} \left( zD_v, z\frac{\partial}{\partial t^0} \right).
\]

Here we use the analytic structure sheaf. Note that we have

\[
D_u = -dD_v - \lambda \frac{\partial}{\partial t^0}, \quad \frac{\partial}{\partial s^0} = \frac{\partial}{\partial t^0}
\]

under the co-ordinate change \((t^0, v) \mapsto (s^0, u)\). Let \(\text{pr} : \hat{\mathcal{M}} \times \mathbf{C}_\zeta \times \mathbf{C}_\lambda \to \mathcal{M}^0 \times \mathbf{C}_\zeta \times \mathbf{C}\) denote the natural projection. The pull back \(\hat{\mathcal{F}}^{\text{tw}} := \text{pr}^* \mathcal{F}^{\text{tw}}\) has the structure of an \(\hat{\mathcal{R}}^{\text{tw}}\)-module by

\[
\frac{\partial}{\partial t^0} \cdot \Delta_v = \Delta_v, \quad v \in \mathcal{B}.
\]

By a solution of the \(\hat{\mathcal{R}}^{\text{tw}}\)-module \(\hat{\mathcal{F}}^{\text{tw}}\), we mean an \(\hat{\mathcal{R}}^{\text{tw}}\)-module homomorphism \(\varphi : \hat{\mathcal{F}}^{\text{tw}}|_V \to \mathcal{O}_V\) for an open subset \(V \subset \hat{\mathcal{M}} \times \mathbf{C}_\zeta \times \mathbf{C}_\lambda.\) We construct a vector-valued solution with values in \(\overline{H}\) such that all of its components form a basis of solutions.

**Definition 5.8.** — The generalized twisted I-functions \(I^{\text{tw},\nu}, \nu \in \mathcal{B}\) are defined as follows (the relevant convergence will be shown in Lemma 5.10 below):

(i) In the FJRW side:

\[
I^{\text{FJRW}}_{\text{tw}}(s^0, u, z; \lambda) = z e^{v_0/z} \sum_{k=0+1}^\infty u^k \prod_{j=1}^N \prod_{0 < b < kq_j + v_j, (b) = (kq_j)} (-q_j \lambda - bz) \phi_{k-1}
\]

where we use the convention of reducing the index \(k - 1\) of \(\phi_{k-1}\) modulo \(d.\) This is an \(H_{\text{ext}}\) valued power series convergent on the region \(|u| < v_c^{-1/d}\) \(\times \mathbf{C}_\zeta^\circ \times \mathbf{C}_\lambda in \hat{\mathcal{M}} \times \mathbf{C}_\zeta \times \mathbf{C}_\lambda.\)

Note that, if \(k \geq v_0 + 1,\) then \(kq_j + v_j \geq q_j + (q_j v_0 + v_j) \geq q_j.\)

(ii) In the GW side (cf. [43, Definition 4.5]):

\[
I^{\text{GW},\nu}(t^0, u, z; \lambda) = z e^{(t^0 + p t^0)/z} \sum_{n \in \mathcal{Q}(u) \in \mathfrak{S}} u^n \prod_{b=1}^d (dp + \lambda + bz) \times \prod_{i=1}^N \prod_{b \leq 0, (b) = (w_{ij})} (w_{ij} p + bz) \prod_{b \leq \nu_i - v_i, (b) = (w_{ij})} (w_{ij} p + bz)^{1-\eta}.
\]
This is an $H_{CR}(\mathcal{P}(w))$-valued power series convergent on the region $\{|v| < v_c\} \times C_{\hat{\varepsilon}}^\times \times C_{\lambda}$ in $\hat{\mathcal{M}} \times C_{\varepsilon} \times C_{\lambda}$. Note that the term

$$\prod_{j=1}^{N} \frac{\prod_{b \leq 0, (b) = (w, p)} (w_{i}p + bz)}{\prod_{b \leq w, n - v_i, (b) = (w, p)} (w_{i}p + bz)} 1_{\{n\}}$$

vanishes if $w, n - v_i < 0$ for all $i$ such that $w, n \in \mathbb{Z}$ (because the factor $\prod_{(w, n) = 0} (w, p)$ in the numerator vanishes on $\mathcal{P}(w)$ for dimensional reason). Thus one can assume that there exists $i$ such that $w, n \in \mathbb{Z}$ and $w, n - v_i \geq 0$. In this case we have $q_i (dn + v_0) + v_i - w, n = q_i v_0 + v_i \geq 0$. Hence one can assume $dn + v_0 \geq 0$ in the summation.

**Remark 5.9.** — For $v = 0$, $I^{w, 0}_{FJRW}(0, u, z)$ and $I^{w, 0}_{GW}(0, v, z)$ coincide with the original twisted I-functions in Section 4.3.1. Also note that the generalized I-function $I^{w, v}$ is homogeneous of degree $2 + \deg \Delta_v = 2(1 + v_0 + \cdots + v_N)$ with respect to the degree $\deg v^0 = \deg \hat{\varepsilon} = \deg z = \deg \lambda = 2$, $\deg u = \deg v = 0$ and the grading on $\hat{\mathcal{H}}$.

**Lemma 5.10.** — The function $I^{w, v}(0, u, z; \lambda)$ is convergent on the region $\{|u| < v_c^{-1/d}\} \times C_{\hat{\varepsilon}}^\times \times C_{\lambda}$ in $\hat{\mathcal{M}} \times C_{\varepsilon} \times C_{\lambda}$; $I^{w, v}(0, v, z; \lambda)$ is convergent on the region $\{|v| < v_c\} \times C_{\varepsilon}^\times \times C_{\lambda}$ in $\hat{\mathcal{M}} \times C_{\varepsilon} \times C_{\lambda}$.

**Proof.** — Write $I^{w, v}_{FJRW}(s^0, u, z; \lambda) = z e^{0/\varepsilon} = z e^{0/\varepsilon} \sum_{k = v_0 + 1}^{\infty} u^k \Delta_k(z, \lambda)$ with $\Delta_k(z, \lambda)$ an element of $H_{ext}$. Fix a norm $\| \cdot \|$ on $H_{ext}$ such that $\|\phi_k\| = 1$. Then we have

$$\frac{\|\Delta_k(z, \lambda)\|}{\|\Delta_0(z, \lambda)\|} = \prod_{j=1}^{N} \prod_{a=0}^{w_j} \frac{|q_{j} \lambda + (k q_{j} + v_j + a)z|}{|k - v_0 + a|}.$$ 

This converges to $\prod_{j=1}^{N} q_{j}^{w_j} = v_c$ as $k \to \infty$ for a fixed $(z, \lambda) \in C_{\hat{\varepsilon}}^\times \times C$. This implies that the convergence radius of $I^{w, v}_{FJRW}$ as a power series of $u$ is $v_c^{-1/d}$. Similarly, write $I^{w, v}_{GW}(0, v, z; \lambda) = z e^{(p + p \log v)/\varepsilon} \sum_{n \in \mathbb{Q}_{>0} \times \mathbb{Z}} \Delta_n(z, \lambda)$ with $\Delta_n(z, \lambda)$ an element of $H_{CR}(\mathcal{P}(w))$. Fix a norm $\| \cdot \|$ on $H_{CR}(\mathcal{P}(w))$. We have

$$\frac{\|\Delta_n(z, \lambda)\|}{\|\Delta_0(z, \lambda)\|} \leq \frac{\prod_{j=1}^{N} \prod_{a=0}^{w_j} (dp + \lambda + (dn + v_0 + a)z)}{\prod_{j=1}^{N} \prod_{a=0}^{w_j} (w_jp + (w_jn - v_j + a)z)}$$

where in the right-hand side $\| \cdot \|$ means the operator norm. The right-hand side converges to $d^d / \prod_{j=1}^{N} w_j^{w_j} = v_c^{-1}$ as $n \to \infty$ for a fixed $(z, \lambda) \in C_{\hat{\varepsilon}}^\times \times C$. Hence the convergence radius of $I^{w, v}_{GW}(0, v, z; \lambda)$ as a power series of $v$ is no less than $v_c$. \hfill \Box

**Proposition 5.11.** — For each $\varphi \in Hom(\hat{\mathcal{H}}, C)$, the map

$$\Gamma^\varphi : \hat{\mathcal{F}}^{sw} \longrightarrow \mathcal{O}, \quad \Delta_v \longrightarrow z^{-1} \varphi(\Gamma^{w, v}), \quad v \in B$$
defines a solution to the $\hat{R}^{tw}$-module $\hat{F}^{tw}$, i.e. a homomorphism of $\hat{R}^{tw}$-modules. Moreover, for a $C$-basis $\varphi_1, \ldots, \varphi_d$ of $\text{Hom}(\hat{H}, C)$, the corresponding solutions $I^{\varphi_1}, \ldots, I^{\varphi_d}$ are linearly independent. (In fact, they form a basis of solutions by Theorem 5.7.)

**Proof.** — For the former statement, it suffices to check that $I^{tw,v} = I^{tw,v}_{\text{FJRW}}$ or $I^{tw,v}_{\text{GW}}$ satisfies the following differential equations (cf. (68); note also the co-ordinate change (70)):

$$
\begin{align*}
(-zD_u + (v_0 + 1)z)I^{tw,v} &= I^{tw,v+0}, \\
(-q_i zD_u - q_i \lambda - v_i z)I^{tw,v} &= I^{tw,v+e_i},
\end{align*}
$$

(71)

Moreover, for a $C$-basis $\varphi_1, \ldots, \varphi_d$ of $\text{Hom}(H, C)$, the corresponding solutions $I^{\varphi_1}, \ldots, I^{\varphi_d}$ are linearly independent. (In fact, they form a basis of solutions by Theorem 5.7.)

They follow from a straightforward computation. Let $\nu(l)$ be as in Lemma 5.4. For the FJRW I-functions, we have

$$
\begin{align*}
z^{-1}I^{tw,v(l)}_{\text{FJRW}} &\sim e^{\delta/(l+1)}(\phi_l + O(u)), \quad l = 0, \ldots, d-1.
\end{align*}
$$

(72)

Since the leading terms span $H_{\text{ext}}$, it follows that the solutions $I^{\varphi_{l1}}_{\text{FJRW}}, \ldots, I^{\varphi_{ld}}_{\text{FJRW}}$ are linearly independent. For the GW I-functions, we have if $\langle l/d \rangle \in \mathfrak{F}$,

$$
\begin{align*}
z^{-1}I^{tw,v(l)}_{\text{GW}} &\sim e^{\delta/(l+d)\log v}/z^{1/d}(1_{\langle l/d \rangle} + O(v^{1/d})).
\end{align*}
$$

(73)

These leading terms span $H_{\text{cr}}(P(w))$. Hence $I^{\varphi_{l1}}_{\text{GW}}, \ldots, I^{\varphi_{ld}}_{\text{GW}}$ are linearly independent. □

The twisted I-function $I^{tw,0}$ in each theory has the $z^{-1}$-asymptotics of the form (cf. (41)):

$$
I^{tw,0} = zF \cdot T_0 + G + O(z^{-1})
$$

(74)

where $F$ and $G$ are functions on a domain in $\hat{M} \times C_\lambda$ on which $I^{tw,0}$ converges; $F$ takes values in $C$ and $G$ takes values in the degree $\leq 2$ part $\hat{H}^{\leq 2} = \hat{H}^0 \oplus \hat{H}^2$. More precisely, in the FJRW side, we have (cf. (42)):

$$
F = F_{\text{FJRW}}(u) = \sum_{k \geq 1, k \equiv 1 (d)} u^k \prod_{j=1}^N (kq_j - 1)_{[kq_j - 1]}(k - 1)!
$$

$$
G = G_{\text{FJRW}}(s^0, u, \lambda)
$$
\[
\begin{align*}
\text{FJRW mirror map} & = s^0 F_{\text{FJRW}}(u) \phi_0 - \sum_{k \geq 2, \sum_{j=1}^{N} (q_j(k-1)) = 1} u^k \frac{\prod_{j=1}^{N} (kq_j - 1)_{[kq_j]-1}}{(k-1)!} \phi_k \\
\text{GW side} & = \lambda \sum_{k \geq d+1 \leq 1(d)} u^k \left( \sum_{i=1}^{N} \sum_{0 < a < q_i, (a) = (kq_i)} \frac{q_i}{a} \prod_{j=1}^{N} (kq_j - 1)_{[kq_j]-1} \right) \frac{1}{(k-1)!} \phi_0
\end{align*}
\]

and in the GW side, we have:

\[
\begin{align*}
F & = F_{\text{GW}}(v) = \sum_{n=0}^{\infty} v^n \frac{(dn)!}{\prod_{j=1}^{N} (w_j n)!} \\
G & = G_{\text{GW}}(\ell^0, v; \lambda) \\
& = \ell^0 F_{\text{GW}}(v) 1 + \sum_{n \in \mathbb{Q}_{>0}, (n) \in \mathbb{R}, \sum_{j=1}^{N} (-w_j n) = 1} v^n \frac{(dn)!}{\prod_{j=1}^{N} (w_j n)!} 1_{(-n)} \\
& + \left[ F_{\text{GW}}(v) \log v + \sum_{n=1}^{\infty} v^n \left( \sum_{a=1}^{d} \frac{d}{a} - \sum_{i=1}^{N} \sum_{a=1}^{w_i n} \frac{w_i}{a} \right) \frac{(dn)!}{\prod_{j=1}^{N} (w_j n)!} \right] \rho 1 \\
& + \lambda \sum_{n=1}^{\infty} v^n \left( \sum_{a=1}^{d} \frac{1}{a} \right) \frac{(dn)!}{\prod_{j=1}^{N} (w_j n)!} 1
\end{align*}
\]

where \((a)_n = a(a-1) \cdots (a-n+1) = \Gamma(a+n)/\Gamma(a+1)\) denotes the falling factorial.

We define the mirror map \(\varsigma\) to be the \(\overline{H}_{\leq 2}\)-valued function:

\[
(75) \quad \varsigma = \frac{G}{F}.
\]

The FJRW mirror map \(\varsigma_{\text{FJRW}}\) is defined over \(\{|u| < v_c^{-1/d}\} \times \mathbb{C}_\lambda\) and the GW mirror map \(\varsigma_{\text{GW}}\) is defined over \(\{|v| < v_c\} \times \mathbb{C}_\lambda\). The following mirror theorem gives a refinement of Theorem 3.10 and [17, 19], namely, the special case \(v = 0\) corresponds to the original mirror theorem (see (78) in the proof). A similar refinement was given in [43, Theorem 4.6] for GW theory of complete intersections in toric orbifolds.

**Theorem 5.12.** — In both FJRW and GW theories, there exist \(\overline{H}\)-valued complex analytic functions \(\Upsilon^{\text{tw}, v}, v \in B\) defined on an open set \(\overline{U} \times \mathbb{C}_z \times \mathbb{C}_\lambda \subset \overline{\mathcal{M}} \times \mathbb{C}_z \times \mathbb{C}_\lambda\) such that

\[
L^{\text{tw}}(\varsigma(x; \lambda), z; \lambda) \Upsilon^{\text{tw}, v}(x, z; \lambda) = z \Upsilon^{\text{tw}, v}(x, z; \lambda).
\]

Here \(\varsigma\) denotes the mirror map (75) in each theory. The fundamental solution \(L(\varsigma(x; \lambda), z; \lambda)\) is also analytic over \(\overline{U} \times \mathbb{C}_z \times \mathbb{C}_\lambda\). The open subset \(\overline{U} \subset \overline{\mathcal{M}}\) is of the form \(\{|u| < \epsilon\}\) for FJRW theory and is of the form \(\{|v| < \epsilon\}\) for GW theory. For \(v = 0\), we have \(\Upsilon^{\text{tw}, 0} = F \cdot T_0\) where \(F\) is the function appearing in (74).
Proof. — First we discuss the case of FJRW theory. By Theorem 3.10, we have

\begin{equation}
F(u)J^\text{tw}(\zeta(s^0, u; \lambda), z; \lambda) = \Gamma^\text{tw,0}(s^0, u, z; \lambda) \quad \text{for } s^0 = 0
\end{equation}

where the subscripts “FJRW” are omitted. We have \( \zeta(s^0, u; \lambda) = s^0T_0 + \zeta(0, u; \lambda) \) and \( \Gamma^\text{tw,0}(s^0, z; \lambda) = e^{0/z}\Gamma^\text{tw}(0, z; \lambda) \). By the string equation for the twisted invariants (see Section 3.2.2), we have

\[ J^\text{tw}(s^0T_0 + \zeta(0, u; \lambda), z; \lambda) = e^{0/z}J^\text{tw}(\zeta(0, u; \lambda), z; \lambda). \]

Hence (77) holds for arbitrary \( s^0 \). Therefore, by Proposition 3.15, we have

\begin{equation}
L^\text{tw}(\zeta(s^0, u; \lambda), z; \lambda)\Gamma^\text{tw,0}(s^0, u, z; \lambda) = zF(u)T_0.
\end{equation}

This shows that one can take \( \Upsilon^\text{tw,0}(s^0, u, z; \lambda) = F(u)T_0 \). The other \( \Upsilon^\text{tw,v} \)'s are obtained from this by differentiation. To see this, we use the fact that the generalized I-functions satisfy (71) and that we have by Proposition 3.12

\[ (\zeta^*\nabla_{\text{D}_u}) \circ L^\text{tw}(\zeta(s^0, u; \lambda), z; \lambda) = L^\text{tw}(\zeta(s^0, u; \lambda), z; \lambda) \circ z\text{D}_u, \]

where \( \zeta^*\nabla_{\text{D}_u} = z\text{D}_u + (D_u\zeta(s^0, u; \lambda)) \cdot \text{tw} \). For example, one obtains \( z\Upsilon^\text{tw,v} \) as

\begin{align*}
L^\text{tw}(\zeta(s^0, u; \lambda), z; \lambda)\Gamma^\text{tw,v}(s^0, u, z; \lambda) &= L^\text{tw}(\zeta(t^0, u; \lambda), z; \lambda)(-q_i z\text{D}_u - q_i\lambda) \cdot \Gamma^\text{tw,0}(s^0, u, z; \lambda) \\
&= (-q_i \zeta^*\nabla_{\text{D}_u} - q_i\lambda)(zF(u)T_0).
\end{align*}

To obtain \( \Upsilon^\text{tw,v} \) for a general \( v \), we use the following differential operator:

\[ P_v(z\text{D}_u) = u^v \prod_{b=1}^{v_0+kd}(-z\text{D}_u + bz) \cdot \prod_{b=-\infty}^{v_0-1}(-q_i z\text{D}_u - q_i\lambda - bz) \]

where \( k \) is an integer such that \( v_0 + kd \geq 0 \). When \( v_i - kw_i < 0 \) for some \( i \), we expand the factor \( (-q_i z\text{D}_u - q_i\lambda - bz)^{-1} \) in the \( \lambda^{-1} \)-series

\[ \sum_{a=0}^{\infty}(-q_i\lambda)^{-a-1}(bz + q_i z\text{D}_u)^a. \]

Then we have \( P_v(z\text{D}_u)\Gamma^\text{tw,0} = \Gamma^\text{tw,v} \). By applying \( P_v(z\text{D}_u) \) to (78), one obtains (76) with \( \Upsilon^\text{tw,v}(s^0, u, \lambda) = P_v(z\text{D}_u)F(u)T_0 \). Note that this expression makes sense as an element of \( \overline{H} \otimes \mathbb{C}[z](\lambda^{-1})([s^0, u]) \). This is because \( \zeta^*\nabla_{\text{D}_u} = z\text{D}_u + (D_u\zeta) \cdot \text{tw} = z\text{D}_u + O(u) \) (note that \( \zeta(s^0, u) = s^0\phi_0 - u\phi_1 + O(u^2) \)). A posteriori, we know that \( \Upsilon^\text{tw,v} \) belongs to \( \overline{H} \otimes \mathbb{C}[z, \lambda][[s^0, u]] \) by (76) since \( L^\text{tw}, \Gamma^\text{tw} \) and \( \zeta \) are regular at \( \lambda = 0 \).
Next we show the analyticity of $\Upsilon^{\text{tw},v}$ and $L^{\text{tw}}$. It suffices to show the analyticity of $L^{\text{tw}}$. By (76) we have

\[
    z^{-1} \begin{bmatrix}
        u^{-1} \Upsilon^{\text{tw},v(0)} & \cdots & u^{-d} \Upsilon^{\text{tw},v(d-1)}
    \end{bmatrix}
\]

\[
    = (L^{\text{tw}})^{-1} \begin{bmatrix}
        u^{-1} \Upsilon^{\text{tw},v(0)} & \cdots & u^{-d} \Upsilon^{\text{tw},v(d-1)}
    \end{bmatrix}
\]

where $v(l)$ is as in Lemma 5.4. Using the basis $\phi_0, \ldots, \phi_{d-1}$ of $\overline{H}$, one can view this as an equality of $(d, d)$ matrices. Write $\gamma(z)$ for the left-hand side. It is invertible near $u = 0$ because of the asymptotics (72). Hence $S^1 \ni z \mapsto \gamma(z)$ defines an element of the loop group $LGL_d = C^\infty(S^1, GL_d)$. As observed by Coates-Givental [21] and Guest [32], we can regard (79) as a Birkhoff factorization of $\gamma(z)$ because $(L^{\text{tw}})^{-1} = \text{id} + O(z^{-1})$ and $\Upsilon^{\text{tw},v}$ is regular at $z = 0$. Birkhoff’s theorem [58] says that the multiplication map $LGL_d \times LGL_d \to LGL_d$ is an isomorphism onto an open and dense subset called the “big cell”. Here $LGL_d$ is the subgroup consisting of the boundary values of holomorphic maps $\gamma_- : \{ z \in \mathbb{C} \cup \{ \infty \} \mid |z| < 1 \} \to GL_d$ satisfying $\gamma_-(\infty) = \text{id}$ and $L^+GL_d$ is the subgroup consisting of the boundary values of holomorphic maps $\gamma_+ : \{ z \in \mathbb{C} \mid |z| > 1 \} \to GL_d$. The asymptotics (72) ensures that $\gamma(z)$ is in the big cell for $|u| < \epsilon$ and $|\lambda| \leq 1$ for sufficiently small $\epsilon > 0$. Hence its Birkhoff factor $\gamma_-(z) = L^{\text{tw}}(z^0, u; \lambda, z; \lambda)^{-1}$ is analytic on the region $\{|u| < \epsilon, |z| > 1, |\lambda| < 1\}$. The homogeneity of $(L^{\text{tw}})^{-1}$ implies that it is in fact analytic on $\{|u| < \epsilon, z \in \mathbb{C}^x, \lambda \in \mathbb{C}\}$.

For GW theory, the theorem follows from the proof of [43, Theorem 4.6]. When $(l/d) \in \mathfrak{F}$, the function $\Upsilon_{GW}^{\text{tw},v(l)}$ coincides with $zv^{-l/d} \Upsilon_{\mathfrak{F}}^{\text{tw},v(l)}|_{Q=1}$ there and we can take $\Upsilon_{\mathfrak{F}}^{\text{tw},v(l)} = v^{l/d} \Upsilon_{\mathfrak{F}}^{\text{tw},v(l)}|_{Q=1}$ in the notation of loc. cit. We can get the other $\Upsilon_{GW}^{\text{tw},v(l)}$ from $\Upsilon_{GW}^{\text{tw},v(l)}$, $(l/d) \in \mathfrak{F}$ by applying differential operators in $\widehat{R}_{\text{tw}}$, so the other $\Upsilon^{\text{tw},v}$ as well. $\square$

Let $\widehat{U}_{\text{FJR}} = \{|u| < \epsilon\}, \widehat{U}_{\text{GW}} = \{|v| < \epsilon\}$ be sufficiently small open subsets of $\widehat{M}$ as in Theorem 5.12. The following corollary gives a twisted version of Theorem 2.23.

**Corollary 5.13.** — Via the $\widehat{R}_{\text{tw}}$-module $\widehat{F}_{\text{tw}}$ over $\widehat{M}$, the $e_{\mathbb{C}^x}$-twisted quantum connections $\nabla^{\text{tw}}$ of FJRW theory and of GW theory are analytically continued to each other. More precisely, we have a local trivialization of $\widehat{F}_{\text{tw}}$ over $\widehat{U}_{\nabla}$

\[
    \text{Mir}_{\nabla} : \widehat{F}_{\text{tw}}|_{\widehat{U}_{\nabla} \times \mathbb{C}^x} \cong \overline{H} \otimes \mathcal{O}_{\widehat{U}_{\nabla} \times \mathbb{C}^x} \quad \nabla = \text{FJRW or GW}
\]

such that, under the trivialization, the action of $\widehat{R}_{\text{tw}}$ is given by the $e_{\mathbb{C}^x}$-twisted quantum connection.
\[ zD \longmapsto \zeta^* \nabla_{zD}, \quad D \text{ is a vector field on } \hat{U}, \]

where \( \zeta^* : \hat{U} \times C_\lambda \rightarrow \overline{H} \) is the mirror map (75) in the respective theory.

**Proof.** — We omit the subscript “FJRW” or “GW” throughout the proof. By Proposition 5.11, the generalized twisted I-functions define an \( \overline{H} \)-valued solution:

\[
(80) \quad \hat{F}\tw \mid_{\hat{U} \times C_\lambda} \rightarrow \overline{H} \times \mathcal{O}_{\hat{U} \times C_\lambda}, \quad \Delta_\nu \mapsto z^{-1} I_{\tw, \nu},
\]

which is an isomorphism (see the asymptotics (72), (73)). On the other hand, the twisted quantum connection \( \zeta^* \nabla_{\tw} \) also has an \( \overline{H} \)-valued solution (Proposition 3.12)

\[
L_{\tw} \left( \zeta(\cdot; \lambda), z; \lambda \right)^{-1} : \left( \overline{H} \otimes \mathcal{O}_{\hat{U} \times C_\lambda}, \zeta^* \nabla_{\tw} \right) \rightarrow \overline{H} \otimes \mathcal{O}_{\hat{U} \times C_\lambda},
\]

which sends \( \Upsilon_{\tw, \nu} \) to \( z^{-1} I_{\tw, \nu} \) by Theorem 5.12. This is also an isomorphism. Therefore we have an isomorphism \( \text{Mir} : \hat{F}\tw \mid_{\hat{U} \times C_\lambda} \cong \overline{H} \otimes \mathcal{O}_{\hat{U} \times C_\lambda} \) such that \( \text{Mir}(\Delta_\nu) = \Upsilon_{\tw, \nu} \). It extends across \( z = 0 \) as \( \Upsilon_{\tw, \nu} \) is regular at \( z = 0 \). Now it suffices to show that \( \Upsilon_{\tw, \nu}, \nu \in B \) generate \( \text{Holl} \) at \( z = 0 \). In the case of FJRW theory, this follows from the fact that the factor \( [u^{-1} \Upsilon_{\tw, \nu}(0), \ldots, u^{-d} \Upsilon_{\tw, \nu}(d-1)] \) in the Birkhoff factorization (79) is invertible at \( z = 0 \). The discussion is similar for GW theory. □

**Remark 5.14.** — Mann-Mignon [50, Theorem 1.2] described explicitly the twisted quantum D-module (with \( \lambda = 0 \)) for a smooth nef complete intersection in a toric manifold.

### 5.3. Analytic continuation \( \mathbf{U}_{\tw} \) revisited.

**Lemma 5.15.** — The submodule \( \mathcal{R}_{\tw} \Delta_0 \) of \( \mathcal{F}_{\tw} \) coincides with \( \mathcal{F}_{\tw} \) at the generic point on \( \mathcal{M} \times C_\zeta \times C_\lambda \).

**Proof.** — In view of the isomorphism (80), it suffices to show that \( \prod_{l=1}^{l} (zD_u - b z) z^{-1} I_{\tw,0}^{FW, \nu} \), \( l = 0, \ldots, d - 1 \) form a basis of \( \overline{H} = \text{Hext} \) for generic \( (z, \lambda) \) and sufficiently small \( |u| \). This follows from the asymptotics near \( u = 0 \):

\[
\prod_{l=1}^{l} (zD_u - b z) z^{-1} I_{\tw,0}^{FW, \nu} \sim e^{b/|u|} u^{l+1} (-1)^l \left( \prod_{j=1}^{N} \prod_{0 < b < (l+1)q_j}^{(l)} \prod_{(\theta) = ((l+1)q_j)}^{(l)} (-q_j \lambda - b z) \right) \phi_l + O(u^{l+2}).
\]

We calculated in Proposition 4.12 a linear transformation \( \mathbf{U}_{\tw} (64) \)

\[
\mathbf{U}_{\tw} : \text{Hext} \otimes \mathcal{O}_{\Delta_\xi} \rightarrow \text{HCR}(\mathbf{P}(\nu)) \otimes \mathcal{O}_{\Delta_\xi}
\]
from analytic continuation of the “$H$-functions”, where $\Delta_\xi = \{ |\xi| < \varepsilon \}$ denotes a sufficiently small disc in the $\xi := (\lambda/z)$-plane. We give an interpretation of $U^\tw$ as analytic continuation of flat sections of the global D-module $\hat{F}^\tw$. Using the trivialization $\text{Mir}_\varpi$ in Corollary 5.13, we define a flat section $f^\tw_\varpi(\alpha)$ of $\hat{F}^\tw$ over $\hat{U}_\varpi \times \{(z, \lambda) \in \mathbb{C}^\times \times \mathbb{C} \mid |\xi| = |\lambda/z| < \varepsilon \}$ parametrized by $\alpha \in \overline{\varpi} \otimes \mathcal{O}(\Delta_\xi)$:

\[
(81) \quad f^\tw_\varpi(\alpha)(x, z; \lambda) := L^\tw_\varpi((\zeta(x; z), z; \lambda))e^{\text{Gr}(\hat{F}^\tw)(2\pi \delta(z, \lambda))}, \quad \alpha \in \overline{\varpi} \otimes \mathcal{O}(\Delta_\xi)
\]

where $L^\tw_\varpi, \hat{\Gamma}^\tw_\varpi$ are the fundamental solution and the twisted Gamma class (Section 4.3.3) in each theory and $\varpi = \text{GW}$ or FJRW. We extend the $H$-functions in the $s^0$- or $l^0$-direction as follows:

\[
\begin{align*}
\hat{H}^\tw_{\text{FJRW}}((s^0, u), z; \lambda) &= e^{s^0/z}\hat{H}^\tw_{\text{FJRW}}(u, z; \lambda) \\
\hat{H}^\tw_{\text{GW}}((l^0, v), z; \lambda) &= e^{l^0/z}\hat{H}^\tw_{\text{GW}}(v, z; \lambda)
\end{align*}
\]

so that we have $I^\tw_\varpi = z^{-\text{Gr}(\hat{F}^\tw)(2\pi \delta(z, \lambda)))}$ (cf. (59)). By this relation and Theorem 5.12, we have

\[
(82) \quad \text{Mir}_\varpi(z\Delta_0) = z\hat{\gamma}^\tw_{\varpi, 0}(x, z; \lambda) = f^\tw_\varpi(\hat{H}^\tw_{\varpi}(x, z; \lambda))(x, z; \lambda).
\]

Namely the $H$-function represents the section $z\Delta_0$ in the flat frame $f^\tw_\varpi$. The relationships between $\hat{\gamma}^\tw_{\varpi, 0}, z\Delta_0, I^\tw_\varpi, \hat{H}^\tw_{\varpi}$ are given in the following diagram:

\[
\begin{align*}
\mathcal{F}^\tw|_{\mathcal{B}_\varpi} &\xrightarrow{\text{Mir}_\varpi} (\overline{\varpi} \otimes \mathcal{O}_{\mathcal{B}_\varpi}, \zeta(x; z)\nabla^\tw) &\xleftarrow{f^\tw_\varpi} (\overline{\varpi} \otimes \mathcal{O}_{\mathcal{B}_\varpi}, \delta) &\xleftarrow{\text{Mir}_\varpi} (\overline{\varpi} \otimes \mathcal{O}_{\mathcal{B}_\varpi}, \delta) \xrightarrow{f^\tw_\varpi} (\overline{\varpi} \otimes \mathcal{O}_{\mathcal{B}_\varpi}, \delta) \\
\hat{\gamma}^\tw_{\varpi} &\xrightarrow{z\Delta_0} \hat{\gamma}^\tw_{\varpi, 0} &\xleftarrow{I^\tw_\varpi} \hat{H}^\tw_{\varpi}
\end{align*}
\]

where $\mathcal{B}_\varpi = \hat{U}_\varpi \times \{(z, \lambda) \in \mathbb{C}^\times \times \mathbb{C} \mid |\lambda/z| < \varepsilon \}$ and $\delta$ stands for the trivial connection (all the connections here are only defined in the $\hat{U}_\varpi$-direction).

Recall the path $\gamma_l$ in $\hat{M}$ (Section 4.3.4, Figure 3) defined for each integer $l$. It can be lifted to a path $\hat{\gamma}_l$ in $\hat{M}$ starting from the GW base point $\log v \ll 0, l^0 = 0$ and ending at the FJRW base point $\log u \ll 0, s^0 = 0$. The homotopy type of the lift $\hat{\gamma}_l$ is unambiguous. For convenience, we take the following lift $\hat{\gamma}_l$:

\[
\begin{align*}
(\log v \ll 0, l^0 = 0) &\xrightarrow{\sim} (\log u \ll 0, l^0 = 0) = (\log u \ll 0, s^0 = \lambda \log u) \quad \text{(moving along $\gamma_l$)} \\
\log u \ll 0, l^0 = 0 &\xrightarrow{\sim} (\log u \ll 0, s^0 = 0) \quad \text{(shifting $s_0$)}.
\end{align*}
\]
Because the shift of $s_0$ has an effect of the multiplication by the factor $u^{\lambda/z} = u^{\xi}$ on the $\mathcal{H}$-function $\mathcal{H}_{\text{FJRW}}$, we have from Proposition 4.12 that

\begin{equation}
(83) \\
(\mathcal{H}^{\text{GW}})_{\text{continued}} = U^{\text{tw}}_l (\mathcal{H}^{\text{FJRW}}_{\text{tw}})
\end{equation}

where the left-hand side now denotes the analytic continuation of $\mathcal{H}^{\text{GW}}_{\text{tw}}$ along $\hat{\gamma}_l$.

**Proposition 5.16.** — Along the path $\hat{\gamma}_l^{-1}$, the FJRW flat section $\mathcal{H}^{\text{FJRW}}_{\text{tw}} (\alpha)$, $\alpha \in H_{\text{ext}}$ is analytically continued to the GW flat section $\mathcal{H}^{\text{GW}}_l (U^{\text{tw}}_l \alpha)$. Here the values of $z$ and $\lambda$ are fixed during the analytic continuation and chosen so that $|\xi| = |\lambda/z|$ is sufficiently small.

**Proof.** — Note that $\hat{\mathcal{F}}^{\text{tw}}_{\diamond}$ is invertible for sufficiently small $|\xi|$. Therefore $\{\mathcal{H}^{\text{FJRW}}_{\text{tw}} (T_i)\}_{i=0}^{d-1}$ forms a basis of flat sections for sufficiently small $\xi = \lambda/z$. Hence for a fixed such $(\xi, \lambda)$, there exists an invertible linear transformation $V_l : H_{\text{ext}} \to H_{\text{CR}} (P(w))$ such that $\mathcal{H}^{\text{GW}}_l (V_l T_i)$ is analytically continued to $\mathcal{H}^{\text{FJRW}}_{\text{tw}} (V_l T_i)$ along $\hat{\gamma}_l^{-1}$. Because $\mathcal{H}^{\text{GW}}_{\text{tw}} (\mathcal{H}^{\text{GW}}_{\text{tw}}) = \text{Mir}_{\diamond} (z\Delta_0)$ (82) and $z\Delta_0$ is a global section of $\hat{\mathcal{F}}^{\text{tw}}_{\diamond}$, we have

\[ (\mathcal{H}^{\text{GW}}_{\text{tw}})_{\text{continued}} = V_l (\mathcal{H}^{\text{FJRW}}_{\text{tw}}). \]

Because $z\Delta_0$ is a generator of $\hat{\mathcal{M}}\circ \times C_z$ at the generic point (Lemma 5.15), this relation uniquely determines $V_l$ for a generic $(\xi, \lambda)$. By (83), we know that $V_l = U^{\text{tw}}_l$. \qed

**5.4. The non-equivariant limit and its reduction.** — Here we prove Theorem 2.23. By taking the non-equivariant limit $\lambda = 0$ in Corollary 5.13, we obtain analytic continuation between $e$-twisted quantum connections. (Recall that $e$ stands for the non-equivariant Euler class.) We shall show that it reduces to analytic continuation between ambient and narrow part quantum D-modules. This reduction was described more explicitly in terms of the Picard-Fuchs ideal in a recent paper of Mann-Mignon [50, Theorem 1.2] for the quantum cohomology of a smooth nef complete intersection in a toric manifold.

**(Step 0).** — Note that $\hat{\mathcal{M}}\circ \times C_z$ is contained in $\mathcal{M} \times C_z \times C_{\lambda}$ as the locus $\{\lambda = s_0 = 0\} = \{\lambda = s_0 = 0\}$. We consider the restriction

\[ \tilde{\mathcal{G}} := \hat{\mathcal{F}}^{\text{tw}} |_{\lambda = s_0 = 0} \]

of $\hat{\mathcal{F}}^{\text{tw}}$ to $\tilde{\mathcal{M}}\circ \times C_z$. This is also identified with the pull-back of

\[ \mathcal{G} := \mathcal{F}^{\text{tw}} |_{\lambda = 0} \]

by $\tilde{\mathcal{M}}\circ \times C_z \to \mathcal{M} \times C_z$. Let $\tilde{U}_{\diamond}$ denote the open subset of $\tilde{\mathcal{M}}\circ$ given by $\tilde{U}_{\text{GW}} = \{|v| < \epsilon\}$ or $\tilde{U}_{\text{FJRW}} = \{|w| < \epsilon\}$ where $\epsilon$ is the same as in Corollary 5.13. Over $\tilde{U}_{\diamond} \times C_z$, $\mathcal{G}$ is identified with the $e$-twisted quantum connection $\nabla^{\text{tw}}$ on $\overline{H}_{\diamond} \times (\tilde{U}_{\diamond} \times C_z) \to \tilde{U}_{\diamond} \times C_z$. 

by Corollary 5.13. By Proposition 3.11, under the natural projection \( \pr: \overline{H} \to H' \), the \( \epsilon \)-twisted quantum connection projects to the quantum connection of the respective theory:

\[
\begin{align*}
\mathcal{G}|_{\tilde{U} \times_c c} & \overset{\cong}{\longrightarrow} (\overline{H} \otimes \mathcal{O}_{\tilde{U} \times_c c}, \zeta_0^* \nabla_{tw}) \\
\quad \downarrow \pr & \quad \downarrow \pr \\
\quad (H' \otimes \mathcal{O}_{\tilde{U} \times_c c}, (\pr \circ \zeta_0)^* \nabla),
\end{align*}
\]

where \( \zeta_0: \tilde{U} \to \overline{H} \) denotes the mirror map (75) restricted to \( \lambda = \epsilon = 0 \). Here the meromorphic flat connection \( \nabla \) on \( \mathcal{G} \) (or \( \mathcal{G} \)) is given by the action of \( zD_v \in R_{tw}|_{\lambda=0} \), i.e. we define \( \nabla_{D_v} := z^{-1} \) (the action of \( zD_v \)) on \( \mathcal{G} \) (or \( \mathcal{G} \)).

(Step 1). — Let \( U_0 \subset M^0 \) be the image of \( \tilde{U} \) under the projection \( \tilde{M}^0 \to M^0 \). We show that the diagram (84) descends to the quotient \( \overline{U} \to U_0 \). First notice that the Galois symmetry in Propositions 2.14, 2.16 extends to the twisted theory. The map \( G: H' \to H' \) there is extended to \( \overline{H} \) as

\[
G(\phi) = e^{-2\pi i \frac{\lambda}{d}} \phi \quad \text{for FJRW theory;}
\]

\[
G(1) = e^{2\pi i f} (1 - 2\pi i \lambda) \quad \text{for GW theory.}
\]

Then the conclusions of Propositions 2.14, 2.16 hold for this \( G \) (except that we do not have the connection in the \( z \)-direction in the twisted theory). The proof is similar. This shows that the fundamental solution \( L_{tw} \) in the twisted theory (see Proposition 3.12) has the following symmetry:

\[
e^{-2\pi i/d} G \circ L_{FJRW}^{tw}(G^{-1}(t, z; \lambda)) = L_{FJRW}^{tw}(t, z; \lambda) \circ e^{-2\pi i/d} G
\]

\[
dG \circ L_{GW}^{tw}(G^{-1}(t, z; \lambda)) = L_{GW}^{tw}(t, z; \lambda) \circ e^{-2\pi i f/2} \circ dG.
\]

On the other hand, the deck transformation of \( \tilde{U} \) acts on \( L_{tw,v} \) as

\[
e^{-2\pi i/d} G(l_{FJRW}^{tw,v}(\log u + (2\pi i/d), z)|_{\rho=\lambda=0}) = l_{FJRW}^{tw,v}(\log u, z)|_{\rho=\lambda=0}
\]

\[
e^{-2\pi i f/2} \circ dG(l_{GW}^{tw,v}(\log v + 2\pi i, z)|_{\rho=\lambda=0}) = l_{GW}^{tw,v}(\log v, z)|_{\rho=\lambda=0}.
\]

Hence the mirror maps (with \( \epsilon = \lambda = 0 \)) satisfy

\[
G(\zeta_{FJRW}(\log u + (2\pi i/d))) = \zeta_{FJRW}(\log u),
\]

\[
G(\zeta_{GW}(\log v + 2\pi i)) = \zeta_{GW}(\log v).
\]

This shows that the deck transformation of \( \tilde{U} \) is conjugate to the Galois action on \( \overline{H} \) via the mirror maps. By the relation (76) and the above calculations, we find that (again
over the locus $\lambda = \ell^0 = 0$

$$e^{-2\pi i/d} G \left( \bar{T}_{\text{FJRW}}^\text{tw,v} \left( \log u + \frac{2\pi \mathbf{1}}{d}, z \right) \right) = \bar{T}_{\text{FJRW}}^\text{tw,v} \left( \log u, z \right)$$

$$dG \left( \bar{T}_G^\text{tw} \left( \log v + 2\pi \mathbf{1}, z \right) \right) = \bar{T}_G^\text{tw} \left( \log v, z \right).$$

This shows that the induced Galois symmetry on the sheaf $(\bar{H} \times \mathcal{O}_{\bar{U} \times C_z}, \zeta^* \nabla^\text{tw})$ is compatible with the deck transformation on $\tilde{\mathcal{G}}|_{\bar{U} \times C_z}$ because the deck-transformation-invariant section $\Delta_v \in \tilde{\mathcal{G}}$ corresponds to $\nabla^\text{tw}$. Moreover the projection $\text{pr}: \bar{H} \rightarrow H'$ is compatible with the Galois action, so the diagram (84) descends to

$$\mathcal{G}|_{U \times C_z} \xrightarrow{\cong} \left( \bar{H} \times \mathcal{O}_{\bar{U} \times C_z}, \zeta^* \nabla^\text{tw} \right)/\langle G \rangle$$

(86)

$$\downarrow \text{pr}$$

$$(H' \otimes \mathcal{O}_{\bar{U} \times C_z}, (\text{pr} \circ \zeta\circ)^* \nabla)/\langle G \rangle.$$ 

Notice that the bundle in the second line is the pull-back of the quantum D-module $(F, \nabla)/\langle G \rangle$ in Definition 2.10 by the mirror map

$$\tau\circ := [\text{pr} \circ \zeta\circ]: U \rightarrow H'_{2}/\langle G \rangle.$$ 

In the diagram (86), we do not consider the flat connection $\nabla_z$ in the $z$-direction and the pairing $P$. However, we can introduce $\nabla_z$ for $\mathcal{G}$ and make the diagram compatible with $\nabla_z$ as follows. Recall that (the module of global sections of) $F^\text{tw}$ is $2\mathbb{Z}$-graded by $\text{deg } u = \text{deg } v = 0, \text{deg } \Delta_v = 2 \sum_{i=0}^{N} v_i, \text{deg } z = \text{deg } \lambda = 2$. Thus $\mathcal{G} = F^\text{tw}|_{\lambda = 0}$ is also graded. The grading defines the meromorphic flat connection $\nabla_z$ on $\mathcal{G}$ (with logarithmic poles along $z = 0$) as

$$\nabla_z \Delta_v = \frac{1}{z} \frac{\text{deg } \Delta_v}{2} \Delta_v.$$ 

Because all the morphisms in the diagram (86) preserve the grading and the Euler vector field vanishes on the image of the mirror map $\text{pr} \circ \zeta\circ$, the projection $\text{pr}: \mathcal{G}|_{U \times C_z} \rightarrow (\tau\circ)^*(F, \nabla)/\langle G \rangle$ induced from the diagram (86) preserves the connection $\nabla_z$ as well.

(Step 2). — The diagram (86) defines for each $(x, z) \in U \times C_z$ a projection $\mathcal{G}_{(x, z)} \rightarrow H'$, i.e. an element of the Grassmannian $Gr(\mathcal{G}_{(x, z)})$. The kernel of the projection is flat for $\nabla$ (including the $z$-direction). We show that this section of the Grassmannian bundle $Gr(\mathcal{G})$ extends globally over $M^\circ \times C_z$.

Recall the flat section $f^\circ_{\text{pr}}(\alpha)$ of the twisted theory in (81). When restricted to the locus $\lambda = \ell^0 = 0$, this defines a flat section of $\mathcal{G}$. On the other hand, we can define a flat section of the quantum D-module $(H' \otimes \mathcal{O}_{\bar{U} \times C_z}, (\text{pr} \circ \zeta\circ)^* \nabla)$ by an analogous formula:

$$f_{\circ}(\alpha) = L_{\circ} \left( (\text{pr} \circ \zeta\circ)(x), z \right) z^{-\text{deg} \bar{T}_{\circ} \left( (2\pi \mathbf{1}) \frac{\text{deg } \alpha}{2} \right)}, \quad \alpha \in H'_{\circ},$$

(87)
where \( L_\mathcal{G}(t, z) \) and \( \hat{\Gamma}_\mathcal{G} \) are the fundamental solution and the Gamma class in the respective theory (as appear in Definition 2.19). By Proposition 3.14 and the definitions of \( f_\mathcal{G} \) and \( f_\mathcal{C} \), we have

\[
\Pr\{f_\mathcal{G}(\alpha)|_{\lambda=\rho=0}\} = f_\mathcal{C}(\Pr(\alpha))
\]

for \( \alpha \in \mathbb{H} \).

**Lemma 5.17.** — The section of \( \text{Gr}(\mathcal{G}) \) over \( (\mathcal{U}_{\mathcal{GW}} \cup \mathcal{U}_{\mathcal{FJR}}) \times \mathbb{C}_z \) given by the diagram (88) extends to \( ((\mathcal{U}_{\mathcal{GW}} \cup \mathcal{U}_{\mathcal{FJR}}) \times \mathbb{C}_z) \cup (\mathcal{M}^\circ \times \mathbb{C}_z^\circ) \).

**Proof.** — By the flat connection \( \nabla \) on \( \mathcal{G} \), the section of \( \text{Gr}(\mathcal{G}) \) over \( \mathcal{U}_{\mathcal{GW}} \times \mathbb{C}_z^\circ \) can be extended along any path in \( \mathcal{M}^\circ \times \mathbb{C}_z^\circ \). We see that the given section of \( \text{Gr}(\mathcal{G})|_{\mathcal{U}_{\mathcal{GW}} \times \mathbb{C}_z^\circ} \) is analytically continued to the given section of \( \text{Gr}(\mathcal{G})|_{\mathcal{U}_{\mathcal{FJR}} \times \mathbb{C}_z^\circ} \) along the path \( \gamma_l \) in Section 4.3.4. By considering the \( \lambda = 0 \) limit in Proposition 5.16, we know that \( f_{\mathcal{FJR}}^{tw}(\alpha)|_{\lambda=\rho=0} \) is analytically continued to \( f_{\mathcal{GW}}^{tw}(U_l^{tw}\alpha)|_{\lambda=\rho=0} \) along \( \gamma_l^{-1} \), for \( \alpha \in \mathbb{H}_{\mathcal{FJR}} = \mathbb{H}_{\text{ext}} \). By (88), the projections of these flat sections by \( \Pr \) are \( f_{\mathcal{FJR}}(\Pr(\alpha)) \) and \( f_{\mathcal{GW}}(\Pr(\lim_{\lambda \to 0} U_l^{tw}\alpha)) \). Diagrammatically:

\[
\begin{array}{ccc}
\mathcal{U}_{\mathcal{GW}} \times \mathbb{C}_z^\circ & \xrightarrow{\text{analytic continuation}} & \mathcal{U}_{\mathcal{GW}} \times \mathbb{C}_z^\circ \\
\Pr \downarrow & & \Pr \downarrow \\
f_{\mathcal{FJR}}(\Pr(\alpha)) & \xrightarrow{\text{analytic continuation}} & f_{\mathcal{GW}}(\Pr(\lim_{\lambda \to 0} U_l^{tw}\alpha))
\end{array}
\]

Here we used the fact (Corollary 4.15) that there exists a unique operator \( U_l : H_{\mathcal{FJR}}^2 \to H_{\mathcal{GW}}^2 \) such that \( \Pr \circ (\lim_{\lambda \to 0} U_l^{tw}) = U_l \circ \Pr \). The existence of such an operator shows that the sections \( \Pr \) of \( \text{Gr}(\mathcal{G})|_{\mathcal{U}_{\mathcal{GW}} \times \mathbb{C}_z^\circ} \) and \( \text{Gr}(\mathcal{G})|_{\mathcal{U}_{\mathcal{FJR}} \times \mathbb{C}_z^\circ} \) coincide under analytic continuation along \( \gamma_l \). Because this holds for all the paths \( \gamma_l \) with \( l \in \mathbb{Z} \), the conclusion follows. \( \square \)

**Lemma 5.18.** — The section of \( \text{Gr}(\mathcal{G}) \) in the previous lemma extends to \( \mathcal{M}^\circ \times \mathbb{C}_z \).

**Proof.** — The section of \( \text{Gr}(\mathcal{G}) \) here is flat for \( \nabla \) on \( \mathcal{G} \). Therefore, the corresponding element of \( \text{Gr}(\mathcal{G}_{(x,z)}) \) at \( (x,z) \in \mathcal{M}^\circ \times \mathbb{C}_z \) can be represented by a matrix independent of \( z \) when we write it in terms of the homogeneous basis \( z^{-\deg \Delta_v/2} \Delta_v \), \( l = 0, \ldots, d-1 \) of \( \mathcal{G}_{(x,z)} \). Therefore, via the basis \( \Delta_v \), \( v = 0, \ldots, d-1 \), the section \( x \times \mathbb{C}_z \to \text{Gr}(\mathcal{G}_{(x)}) \) can be represented by an algebraic map \( \mathbb{C}_z \to \text{Gr}(\mathcal{G}) \), which extends across \( z = 0 \) by the completeness of \( \text{Gr}(\mathcal{G}) \). This proves the lemma. \( \square \)

(Step 3). — The previous step shows that there exists a projection \( \mathcal{G} \to \mathcal{F} \) to a locally free sheaf \( \mathcal{F} \) over \( \mathcal{M}^\circ \times \mathbb{C}_z \). The sheaf \( \mathcal{F} \) is equipped with a meromorphic flat connection with simple poles along \( z = 0 \).

\[
\nabla : \mathcal{F} \to \mathcal{F}(\mathcal{M}^\circ \times \{0\}) \otimes \Omega_{\mathcal{M}^\circ \times \mathbb{C}_z}^1
\]
Also $\mathcal{F}$ is isomorphic to the pulled-back quantum D-module $(H_{\mathcal{O}} \otimes O_{U_{\mathcal{O}} \times \mathbb{C}_x}, (\mathcal{P} \circ \mathcal{O})^* \mathcal{V})/\langle \mathcal{F} \rangle$ over the open subset $U_{\mathcal{O}} \times \mathbb{C}_x$. In particular, $\mathcal{F}$ extends across the orbifold point $u = 0$ as an orbi-sheaf with flat connection (i.e. $\mu_d$-equivariant flat bundle on a $d$-fold cover). We denote this extension over $\mathcal{M} \times \mathbb{C}_x$ by the same symbol $\mathcal{F}$.

We claim that there is a global $\mathbb{Z}$-local subsystem $F_\mathcal{O}$ of $(\mathcal{F}|_{\mathcal{M} \times \mathbb{C}_x}, \nabla)$ such that it coincides with the $\hat{\Gamma}$-integral structure over $U_{\mathcal{O}} \times \mathbb{C}_x$. By (89), the flat section $f_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}(\alpha)$, $\alpha \in \text{H}_{\text{nat}}(W, \mu_d)$ is analytically continued to $f_{\mathcal{G}_{W}}(U, \alpha)$ along the path $\gamma_{\hat{t}}^{-1}$. Note that $f(\alpha)$ (87) is related to the flat section $\mathcal{S}(\mathcal{E})$ (22) defining the $\hat{\Gamma}$-integral structure by

$$\mathcal{S}(\mathcal{E}) = \frac{1}{(2\pi i)^t} f(\text{inv ch}(\mathcal{E}))(x, z)$$

where $\mathcal{E}$ is an object of $D^b(X_{\mathcal{W}})$ or $MF_{\mu_d}^r(W)$ such that $\text{ch}(\mathcal{E}) \in \mathcal{H}$. Therefore by Theorem 4.17 we know that

(90) $\mathcal{S}_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}(\mathcal{E})$ is analytically continued to $\mathcal{S}_{\mathcal{G}_{W}}(\Phi_i(\mathcal{E}))$ along $\gamma_{\hat{t}}^{-1}$

for $\mathcal{E} \in MF_{\mu_d}^r(W)$ with $\text{ch}(\mathcal{E}) \in \text{H}_{\text{nat}}(W, \mu_d)$. This shows the existence of a global $\mathbb{Z}$-local system and that the analytic continuation along $\gamma_{\hat{t}}^{-1}$ corresponds to the Orlov equivalence $\Phi_i$.

Finally we show that $\mathcal{F}$ admits a global $\nabla$-flat pairing

$$\mathcal{P}: (-)^* \mathcal{F} \otimes \mathcal{F} \to \mathcal{O}_{\mathcal{M} \times \mathbb{C}_x}, \quad \hat{\gamma} = N - 2,$$

which coincides with the pairings $\mathcal{P}_{GW}$, $(-1)^{N-1} \mathcal{P}_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}$ of the quantum D-modules. In order to see that the global pairing exists over $\mathcal{M} \times \mathbb{C}_x$, in view of (90), it suffices to check that

(91) $(-1)^{N-1} \mathcal{P}_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}( (-)^{s}_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}(\mathcal{E}_1), \mathcal{S}_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}(\mathcal{E}_2)) = \mathcal{P}_{GW}( (-)^{s}_{\mathcal{G}_{W}}(\Phi_i(\mathcal{E}_1), \mathcal{S}_{\mathcal{G}_{W}}(\Phi_i(\mathcal{E}_2)))$

for $\mathcal{E}_1, \mathcal{E}_2 \in MF_{\mu_d}^r(W, \mu_d)$ such that $\text{ch}(\mathcal{E}_i) \in \text{H}_{\text{nat}}(W, \mu_d)$. Recall that the pairing between the flat sections $\mathcal{S}(\mathcal{E})$ coincides with the Euler form up to sign (Proposition 2.21). Because the categorical equivalence preserves the Euler pairing $\chi(\mathcal{E}, \mathcal{F}) = \chi(\Phi_i(\mathcal{E}), \Phi_i(\mathcal{F}))$, (91) follows. The global pairing $\mathcal{P}$ over $\mathcal{M} \times \mathbb{C}_x$ extends across $z = 0$ (with zeros of order $\hat{\gamma}$) by Hartog’s principle because it already extends over $U_{\mathcal{O}} \times \mathbb{C}_x$. The non-degeneracy of $\mathcal{P}/(2\pi i \hat{\gamma})$ along $z = 0$ holds for the same reason.

Now the proof of Theorem 2.23 is complete.

Remark 5.19. — We described the global D-module $\mathcal{F}$ as a quotient of $\mathcal{G} = \mathcal{F}_{tw}|_{\lambda=0}$. In [43, Theorem 6.13], with the aid of mirror symmetry, it was described as a submodule of another multi-GKZ system. We can translate this result in our setting as follows. Define the shift map $\mathcal{S}$: $\mathcal{M}^o \times \mathbb{C}_x \times \mathbb{C}_\lambda \to \mathcal{M}^o \times \mathbb{C}_x \times \mathbb{C}_\lambda$ by $\mathcal{S}(x, z, \lambda) = (x, z, \lambda - z)$. Then the map

$$\sigma: \mathcal{F}_{tw} \to \mathcal{S}^* \mathcal{F}_{tw}, \quad \Delta \mapsto \mathcal{S}^* \Delta_{\mathcal{F}_{\mathcal{J}\mathcal{R}\mathcal{W}}}.$$
is a morphism of $R_{\text{tw}}$-modules by the relations (68). This is an isomorphism at the generic point because both $R_{\text{tw}} \Delta_0$ and $R_{\text{tw}} \Delta_{00}$ equal $F_{\text{tw}}$ at the generic point (see Lemma 5.15; the proof there applies also to $R_{\text{tw}} \Delta_{00}$). However, $\sigma$ is not an isomorphism over $\lambda = 0$ and we have $F = \text{Im}(\sigma|_{\lambda = 0}: F_{\text{tw}}|_{\lambda = 0} \rightarrow (S^*F_{\text{tw}})|_{\lambda = 0})$. See also [50].

5.5. Reconstruction of the big quantum D-module. — Here we prove Theorem 2.25. When $X_W$ is a manifold, the orbifold cohomology consists only of untwisted sectors. In particular $H_{\text{amb}}(X_W)$ is spanned $1, p1, \ldots, p^{\dim X_W}$. This allows us to use the reconstruction theorem [36, 41, 47, 59, 60] to obtain the big quantum cohomology from the small one.

More specifically, we apply the reconstruction theorem of a (TE) structure by Hertling-Manin [36, Theorem 2.5] to the global D-module $(F, V)$ over $M$ (which is itself a (TE) structure). For this, one has to check the injectivity condition (IC) and the generation condition (GC) for $(F, V)$. More concretely, (IC) means

$$zD_v \Delta_0|_{z=0} \neq 0,$$

and (GC) means

$$\{(zD_v)^n \Delta_0 \mid n \geq 0\} \text{ generates } F|_{z=0} \text{ over } O_M.$$

We claim that $(zD_v)^n \Delta_0, n = 0, \ldots, \text{rank } F - 1$ is a basis of $F|_{z=0}$ over the open subsets $U_{\text{FRW}}$ and $U_{\text{GW}}$. (Here $U_{\text{FRW}}$ does not contain $u = 0$.) We work over the cyclic cover $\tilde{U}_\emptyset \subset \tilde{M}_0$ of $U_\emptyset$. First observe that we have D-module isomorphisms (cf. (80)):

$$(F, V)|_{\tilde{U}_\emptyset \times C^*} \xrightarrow{\text{Mir}_\emptyset} (H_\emptyset \otimes O_{\tilde{U}_\emptyset \times C^*}, (\text{pr} \circ \zeta_\emptyset)^* V) \xrightarrow{L_1^{-1}} (H_\emptyset \otimes O_{\tilde{U}_\emptyset \times C^*}, d)$$

$$\Delta_v \mapsto \gamma^v := \text{pr}(\gamma_{\text{tw}, v}) \quad \mapsto z^{-1} I_{\emptyset}^v := z^{-1} \text{pr}(I_{\text{tw}, v}|_{\lambda = 0}).$$

Here the first map is induced from the mirror isomorphism in Corollary 5.13 (see also (84)) and the second map is given by the inverse of the fundamental solution $L(t, z)$ in each theory. The relation $L_1^{-1} \gamma_{\emptyset}(x, z) = z^{-1} I_{\emptyset}^v(x, z)$ follows from (76) and Proposition 3.14. Similarly to (79), the two maps $\text{Mir}_\emptyset, L_1^{-1}$ can be viewed as the Birkhoff factors of the composition since $\text{Mir}_\emptyset$ extends regularly to $z = 0$ and $L_1^{-1}$ extends regularly to $z = \infty$. We want to check that $(zD_v)^i \Delta_0, i = 0, \ldots, \text{rank } F - 1$ form a basis. Under the above map, these sections map to

$$(zD_v)^i z^{-1} I_{\text{GW}}^0 = e^{\log v/z} (p' 1 + O(v^{1/d}))$$

over $\tilde{U}_{\text{GW}}$. From these asymptotics, we know that the matrix with the column vectors $(zD_v)^i z^{-1} I_{\text{GW}}^0, i = 0, \ldots, \text{rank } F - 1$ is Birkhoff factorizable (i.e. in the “big cell”) for sufficiently small $|v|$; this means that $(zD_v)^i \Delta_0, i = 0, \ldots, \text{rank } F - 1$ is a basis of $F|_{z=0}$ over $\tilde{U}_{\text{GW}}$. 
Over $\tilde{U}_{\text{FJRW}}$, the calculation is a little more involved. Instead of $(zD_v)^i$, $i = 0, \ldots, \text{rank} F - 1$, we consider the differential operator $P_i$, $i = 0, \ldots, \text{rank} F - 1$ defined inductively by

$$P_0 = u^{-1}, \quad P_i := u^{-\text{ord}}(z\partial_u)P_{i-1}$$

where $\text{ord}, i \in \mathbb{N}$ is determined by $(z\partial_u)P_{i-1}I_{\text{FJRW}}^0 = O(u^{\text{ord}})$. It suffices to show that $P_i\Delta_0$, $i = 0, \ldots, \text{rank} F - 1$ is a basis of $F|_{z=0}$ since $\{P_i\Delta_0\}$ and $\{(zD_v)^i\Delta_0\}$ are related by an invertible matrix along $z = 0$. We have

$$P_i z^{-1}I_{\text{FJRW}}^0 = e_i \frac{\phi_i^{-1}}{z^{l_i-1}} + O(u)$$

where $k_i$ is the $(i+1)$-th smallest element of the set $\mathbb{N} \subset \{1, \ldots, d-1\}$, $e_i \neq 0$ and $l_i := \deg(\phi_i^{-1})/2$. It is not difficult to show that $l_i = i$ when $X_W$ is a manifold. Therefore the matrix having the column vectors $P_i z^{-1}I_{\text{FJRW}}^0$, $i = 0, \ldots, \text{rank} F - 1$ is Birkhoff factorizable for small $|u|$. The claim now follows also over $\tilde{U}_{\text{FJRW}}$.

Because (IC) and (GC) are open conditions, they hold in a Zariski open subset $M'$ of $M$ containing $U_{\text{GW}}$ and $U_{\text{FJRW}}$. At each point $x \in M'$, we have a universal unfolding [36, Definition 2.3] of $(F, \nabla)|_{(M, x) \times \mathbb{C}}$ over the analytic germ $(M, x) \times (\mathbb{C}^{\text{rank} F - 1}, 0) \times \mathbb{C}$. By the universality, they will patch together to form a global (TE) structure $(F^{\text{ext}}, \nabla^{\text{ext}})$ over $M^{\text{ext}} \supset M'$. By [36, Lemma 3.2], the pairing $P$ over $M \times \mathbb{C}$ extends to $M^{\text{ext}} \times \mathbb{C}$, and we have a (TEP) structure $(F^{\text{ext}}, \nabla^{\text{ext}}, P^{\text{ext}})$. The extension of the $\mathbb{Z}$-local system $F_B$ is automatic.

Next we show that $(F^{\text{ext}}, \nabla^{\text{ext}}, P^{\text{ext}})$ coincides with the “big” quantum D-module over a neighbourhood of $U_{\text{FJRW}}$ or $U_{\text{GW}}$. We review the reconstruction of the big FJRW quantum cohomology. Over $U_{\text{FJRW}}$, we already identified $(F, \nabla, P)$ with the quantum D-module over the image of the mirror map $\tau = \text{pr} \circ \zeta$. We take a basis $\{T_i\}_{i=0}^N$ of $\mathbb{H}$ such that $T_0 = \phi_0$, $T_1 = \phi_1$ and write the big quantum product as $T_0 \bullet T_i = \sum_{k=0}^N C^k_i(t)T_k$, where $t = (t^0, \ldots, t^r)$ is the co-ordinates of $\mathbb{H} = H_{\text{inf}}(W, \mu_d)$ dual to $\{T_i\}_{i=0}^N$. Using the frame $\{T_i\}_{i=0}^N$, one can write the connection $\nabla$ of $F|_{U_{\text{FJRW}}}$ as

$$\nabla_{\tau} = \frac{\partial}{\partial u} + \frac{1}{z} \sum_{i=0}^r \frac{\partial \tau^i(u)}{\partial u} (C_{\alpha\beta}^i(\tau(u)))_{\alpha\beta}.$$ 

Here $\tau(u) = \sum_{i=0}^r \tau^i(u)T_i$ denotes the mirror map. The structure constants $C^k_i(t)$ are a priori formal power series in $t$, but we know from the mirror theorem that the above connection $\nabla_{\tau}$ is convergent. Because $\tau(u) = -u\phi_1 + O(u^2)$, we can use $(u, t^0, t^2, \ldots, t^r) \mapsto \tau(u) + \sum_{j \neq 1} v_j T_j$ as a co-ordinate patch of $\mathbb{H}$ near the origin. We want to reconstruct the connection operators

$$\nabla^{\text{ext}}_{\tau} = \frac{\partial}{\partial u} + \frac{1}{z} A(u, t),$$

where $A(u, t)$ is a convergent power series in $u$, $t$. The structure constants can be uniquely determined by comparing with $\nabla_{\tau}$.
\[ \nabla^\text{ext} = \frac{\partial}{\partial t_i} + \frac{1}{z} C^{\beta}_{\alpha i} \left( \tau(u) + \sum_{j \neq 1} t_j T_j \right), \quad i \neq 1 \]
satisfying \( \nabla^\text{ext} u|_{t=0} = \nabla u, \) \( [\nabla^\text{ext}, \nabla^\text{ext}] = [\nabla^\text{ext}, \nabla^\text{ext}] = 0 \) and \( \nabla_i T_0 = T_i. \) Following the method of [36, Lemma 2.9], [41, Section 4.4], one can solve for such \( C^{\beta}_{\alpha i}(\tau(u) + \sum_{j \neq 1} t_j T_j) \) uniquely as a power series in \( t \). This is because \( T_0 = \phi_0 \) is asymptotic to \( u^{-1} \Delta_0 \) as \( u \to 0 \), so is also a cyclic vector of the action of \( [z \nabla u]|_{z=0} \) for a sufficiently small \( u \neq 0 \).

This reconstruction can be done either over the formal Laurent series ring \( \mathbb{C}((u)) \) or for a fixed small \( u \neq 0 \). In the former case, we recover the big quantum product as a formal power series in \( (u, t) \); in the latter case, we get \( C^{\beta}_{\alpha i}(\tau(u) + \sum_{j \neq 1} t_j T_j) \) as a convergent power series of \( t \) ([36, Lemma 2.9]). Therefore \( C^{\beta}_{\alpha i}(\tau(u) + \sum_{j \neq 1} t_j T_j) \) is a formal power series in \( t \) whose coefficients are analytic functions on \( \{ u \in \mathbb{C} | |u| < \epsilon \} \). Moreover for each \( u \) with \( 0 < |u| < \epsilon \), it is convergent as a power series in \( t \). By [40, Lemma 6.5], such a function is holomorphic in a neighbourhood of \( (u, t) = (0, 0) \). This shows the convergence of the big quantum product and that \( (\mathcal{F}^\text{ext}, \nabla^\text{ext}) \) is isomorphic to the big quantum D-module in a neighbourhood of \( \mathbb{U}_{\text{FJRW}} \). The discussion on the GW side is similar and omitted.

### 5.6. Monodromy and autoequivalences.

Here we prove Theorem 2.26. We study the relationships between monodromy of the global quantum D-module \( \mathcal{F} \) and category equivalences.

An object \( E \) of \( D^b(X_W) \) is said to be spherical [63, Definition 1.1] if \( \text{Hom}^n(E, E) = \text{Hom}(E, E[n]) \) is isomorphic to the cohomology of a sphere, i.e.

\[
\text{Hom}^n(E, E) = \begin{cases} \mathbb{C} & n = 0 \text{ or } \dim X_W \\ 0 & \text{otherwise}. \end{cases}
\]

Seidel-Thomas [63] introduced a functor \( T_E : D^b(X_W) \to D^b(X_W) \), called spherical twist, for a spherical object \( E \). This gives an auto-equivalence with the following property:

\[ T_E(F) \cong \text{Cone}(\text{Hom}^*(E, F) \otimes E \to F). \]

**Example 5.20.** — A line bundle \( \mathcal{O}(i) \) on \( X_W \) is a spherical object (since \( X_W \) is Calabi-Yau).

By Proposition 2.21 and (85), the monodromy of flat sections \( s(\mathcal{E}) \) around the paths \( \gamma_{\text{CY}}, \gamma_{\text{LG}} \) (Figure 2) comes from the autoequivalences \( \mathcal{O}(-1), (1) \) of \( D^b(X_W) \) and \( \text{MF}^{\text{gr}}_{\mathcal{P}_E}(W) \) respectively. We already saw in (90) that the analytic continuation along \( \gamma_{\text{con}}^{-1} \) (Figure 3) is induced by the Orlov equivalence \( \Phi_l \). Thus the monodromy along \( \gamma_{\text{con}}^{-1} \) corresponds to the composition \( \Phi_0 \circ \Phi_{-1} \). The following proposition shows that the monodromy around \( \gamma_{\text{con}}^{-1} \) comes from the spherical twist \( T_0 \).
Proposition 5.21. — For $E \in D^b(X_W)$ such that $\text{ch}(E) \in \text{H}_{\text{amb}}(X_W)$, we have $[\Phi_i\Phi_{i+1}(E)] = [T_{\mathcal{O}(l)}(E)]$ in the numerical $K$-group.

Proof. — Let $[a, b]_q$ be the graded Koszul matrix factorization in Example 4.5. Recall that $\text{ch}([a, b]_q), q \in \mathbb{Z}$ span $H_{\text{nan}}(W, \mu_d)$. Hence by Theorem 4.17, $\text{ch}(\Phi_i([a, b]_q)), q \in \mathbb{Z}$ also span $H_{\text{amb}}(X_W)$ since $U_r: H_{\text{nan}}(W, \mu_d) \cong H_{\text{amb}}(X_W)$. Therefore, it suffices to check that $[T_{\mathcal{O}(l)}\Phi_{i+1}([a, b]_q)] = [\Phi_i([a, b]_q)]$ in the $K$-group. By Proposition 4.11, we have

$$[\Phi_{i+1}([a, b]_q)] = \sum_{j_1 < \cdots < j_r \sum_{i=1}^r w_{j_1} \leq m'} (-1)^{r+1} \left[ \mathcal{O} \left( l + 1 + m' - \sum_{a=1}^r w_{j_a} \right) \right]$$

$$[\Phi_i([a, b]_q)] = \sum_{j_1 < \cdots < j_r \sum_{i=1}^r w_{j_1} \leq m} (-1)^{r+1} \left[ \mathcal{O} \left( l + m - \sum_{a=1}^r w_{j_a} \right) \right]$$

where $m$ (resp. $m'$) is the remainder of $q - l$ (resp. $q - l - 1$) divided by $d$. Because $[T_{\mathcal{O}(l)}E] = [E] - \chi(E(-l))[\mathcal{O}(l)]$, we have

$$[T_{\mathcal{O}(l)}\Phi_{i+1}([a, b]_q)] = \sum_{j_1 < \cdots < j_r \sum_{i=1}^r w_{j_1} \leq m'} (-1)^{r+1} \left[ \mathcal{O} \left( l + 1 + m' - \sum_{a=1}^r w_{j_a} \right) \right]$$

$$+ \sum_{j_1 < \cdots < j_r \sum_{i=1}^r w_{j_1} \leq m'} (-1)^r \chi \left( \mathcal{O} \left( 1 + m' - \sum_{a=1}^r w_{j_a} \right) \right) [\mathcal{O}(l)].$$

Here we use the following fact: For $1 \leq i \leq d$, we have

$$\chi(\mathcal{O}(i)) = \dim H^0(X_W, \mathcal{O}(i)) = \begin{cases} \#\{k_1 \leq \cdots \leq k_i \mid s \geq 0, \sum_{k=1}^i w_{k_i} = i\} & \text{if } 1 \leq i \leq d - 1 \\ \#\{k_1 \leq \cdots \leq k_i \mid s \geq 0, \sum_{k=1}^i w_{k_i} = i\} - 1 & \text{if } i = d. \end{cases}$$

Therefore the second term of the right-hand side of (92) gives

$$[\mathcal{O}(l)](-\delta_{m',d-1} + \sum_{j_1 < \cdots < j_r, k_1 \leq \cdots \leq k_i \sum_{i=1}^r w_{j_1} + \sum_{k=1}^i w_{k_i} = m' + 1, \sum_{i=1}^r w_{j_1} \leq m'} (-1)^r).$$

We claim that for $m' \geq 0$

$$\sum_{j_1 < \cdots < j_r, k_1 \leq \cdots \leq k_i \sum_{i=1}^r w_{j_1} + \sum_{k=1}^i w_{k_i} = m' + 1} (-1)^r = 0.$$
The claim follows from the comparison of the coefficient of $t^{m'+1}$ in the following equality:

$$1 = \frac{(1 - t^{w_1})(1 - t^{w_2}) \cdots (1 - t^{w_N})}{(1 - t^{w_1})(1 - t^{w_2}) \cdots (1 - t^{w_N})}.$$

$$= \sum_{p, q \geq 0} \left( \sum_{\substack{j_1 < \cdots < j_r \sum_{a=1}^r w_{j_a} = p}} (-1)^j \left( \sum_{\sum_{a=1}^r w_{j_a} = q} t^j \right) \right).$$

By the above claim, (93) can be rewritten as

$$[\mathcal{O}(l)](-\delta_{m', d-1} - \sum_{\substack{j_1 < \cdots < j_r \sum_{a=1}^r w_{j_a} = m'+1}} (-1)^j).$$

This gives the second term of the right-hand side of (92).

First consider the case where $m' < d - 1$. In this case, by the above calculation, $[T_{\mathcal{O}(l)} \Phi_{l+1}([a, b]_q)]$ equals $[\Phi_{l([a, b]_q})]$ because $m = m' + 1$. Next consider the case where $m' = d - 1$. In this case, we have $m = 0$ and

$$[T_{\mathcal{O}(l)} \Phi_{l+1}([a, b]_q)] = \sum_{j_1 < \cdots < j_r} (-1)^{r+1} \left[ \mathcal{O} \left( l + d - \sum_{a=1}^r w_{j_a} \right) \right] - [\mathcal{O}(l)].$$

We know from the Koszul complex $E_d$ (56) that the first term in the right-hand side vanishes. Because $m = 0$ we have $[\Phi_{l([a, b]_q})] = -[\mathcal{O}(l)]$. The conclusion follows. 

\[\square\]

**Remark 5.22.** — We should have an isomorphism of functors $T_{\mathcal{O}(l)} \cong \Phi_{l} \circ \Phi_{l+1}^{-1}$, but this does not seem to be proved in the literature. E. Segal [62, Theorem 3.13] showed a similar (object-wise) relationship in the category of B-branes on the LG model $(K_{\mathbb{P}(w)}, \tilde{W})$ (which should be equivalent to $D^b(X_W)$).

We speculate that the relations in the fundamental groupoid of $M$

$$\gamma_{l+1} = \gamma_{LG} \circ \gamma_l \circ \gamma_{CY},$$

$$\gamma_{con} = \gamma_1^{-1} \circ \gamma_0,$$

$$\gamma_{d}^{LG} = \text{id}$$

should be lifted to category equivalences as

$$\Phi_{l+1}^{-1} \cong (1) \circ \Phi_1^{-1} \circ \mathcal{O}(-1),$$

$$T_{\mathcal{O}}^{-1} \cong \Phi_1 \circ \Phi_0^{-1},$$

$$(d) \cong [2].$$
The second relation is conjectural (see Remark 5.22) but the other two are easy to show. Note that the identity in the fundamental groupoid is lifted to the 2-shift [2] in the third relation. This is the reason why we have to mod out by [2] in the statement of Theorem 2.26.

Finally we check the last statement in Theorem 2.26. The fundamental group of \( M \) is generated by \( \gamma_{\text{CY}} \), \( \gamma_{\text{con}} \) and is defined by the relation

\[
(\gamma_{\text{CY}} \circ \gamma_{\text{con}})^d = \text{id}.
\]

We define the lift \( \hat{\rho} : \pi_1(M, b_0) \to \text{Auteq}(D^b(X))/[2] \) by sending \( \gamma_{\text{CY}} \) to \( O(-1) \) and \( \gamma_{\text{con}} \) to \( T^{-1}O \) as we speculated above. It suffices to check the relation:

\[
\left(O(-1) \circ T^{-1}O\right)^d \cong [2].
\]

This was proved by Canonaco-Karp [9]. The proof of Theorem 2.26 (hence of Theorem 1.2) is now complete.

\section*{Acknowledgements}

We would like to thank the Institut Fourier, where a major portion of this work was done. A.C. is supported by the ANR grant “Nouvelles symétries pour la théorie de Gromov-Witten”, ANR-09-JCJC-0104-01; he thanks Claire Voisin for a motivating question on the possibility of a direct proof of the LG/CY correspondence without passing through mirror symmetry and José Bertin for dozens of conversations and for his notes on matrix factorizations. H.I. is supported by Grant-in-Aid for Young Scientists (B) 22740042; he thanks Ed Segal for useful discussions on matrix factorizations and Orlov equivalence. Y.R. is supported by a grant from NSF; his two visits to Grenoble were supported by the Institut Fourier and by a “chaire d’excellence ENS Lyon/UJF Grenoble”.

\section*{Appendix A: Proof of Proposition 2.1}

When \( N_k = 0 \), both sides of (7) are one-dimensional and their pairings match. When \( N_k = 1 \), both sides are zero. Assume that \( N_k \geq 2 \). The relative cohomology exact sequence identifies \( H^{N_k}((\mathbb{C}^N)_k, W_k^{+\infty}) \) with \( H^{N_k-1}(W_k^{+\infty}) \cong H^{N_k-1}(W_k^{-1}(1)) \). Therefore

\[
H(W, \mu_d)_k \cong H^{N_k-1}(W_k^{-1}(1))^{\mu_d}.
\]

We use the following result of Steenbrink:
Theorem A.1 ([64, Theorem 1]). — The Deligne weight filtration \( \mathcal{W}_\bullet \) on \( H^{N_i-1}(W^{-1}_k(1)) \) is of the form

\[
0 = \mathcal{W}_{N_i-2} \subset \mathcal{W}_{N_i-1} \subset \mathcal{W}_{N_i} = H^{N_i-1}(W^{-1}_k(1)).
\]

Take a set \( \{ \varphi_1, \ldots, \varphi_k \} \subset C^*_w(W^{-1}_k) \) of homogeneous \( N_k \) forms which gives a basis of \( \Omega(W_k) \). Let \( |i| \) denote the degree of \( \varphi_i \), divided by \( d \). Define \( \eta_i \in H^{N_i-1}(W^{-1}_k(1)) \) by

\[
\eta_i := c_i \text{Res}_{W_k(x) = 1} \left( \frac{\varphi_i}{(W_k(x) - 1)^{|i|}} \right)
\]

with \( c_i = \Gamma(1 - (-|i|))(|i| - 1)! \). Then the set \( \{ \eta_i \mid N_k - 1 - p < |i| < N_k - p \} \) gives a basis of \( \text{Gr}^p_f(\mathcal{W}_{N_i-1}) \); the set \( \{ \eta_i \mid |i| = N_k - p \} \) gives a basis of \( \text{Gr}^p_f(\mathcal{W}_{N_i}/\mathcal{W}_{N_i-1}) \).

There is a typo in the statement of [64, Theorem 1] about the index of the Hodge filtration and we corrected it above. The prefactor \( c_i \) is not important in the above statement, but is chosen for our later purpose. Since \( \{ \eta_i \mid |i| \in \mathbb{Z} \} \) gives a basis of the \( \mu_d \)-invariant part of \( H^{N_i-1}(W^{-1}_k(1)) \), by the theorem, the \( \mu_d \)-invariant part splits the weight filtration:

\[
\mathcal{W}_{N_i}/\mathcal{W}_{N_i-1} \cong H^{N_i-1}(W^{-1}_k(1))^{\mu_d}.
\]

Therefore the sector \( H(W, \mu_d)_k \cong H^{N_i-1}(W^{-1}_k(1))^{\mu_d} \) has a pure Hodge structure of weight \( N_k \). Moreover the theorem gives an isomorphism

\[
\Omega(W_k)^{\mu_d} \cong \text{Gr}^N_f H^{N_i-1}(W^{-1}_k(1))^{\mu_d}, \quad [\varphi_i] \longmapsto [\eta_i]
\]

independent of the choice of representatives \( \varphi_i \). The isomorphism (7) is defined by the Hodge decomposition which splits the above isomorphism:

\[
H(W, \mu_d)_k \cong H^{N_i-1}(W^{-1}_k(1))^{\mu_d} = \bigoplus_{p=0}^{N_i} \mathcal{F}_p \cap \overline{\mathcal{F}}^{N_i-p} \cong \Omega(W_k)^{\mu_d}.
\]

Next we study the pairing on the FJRW state space. The form \( e^{-W_k} \varphi_i \) defines a cohomology class in \( H^{N_i}((\mathcal{C}^N)_k, W^{+\infty}_k) \) via the integration over non-compact Lefschetz thimbles \( \Gamma \in H_{N_i}((\mathcal{C}^N)_k, W^{+\infty}_k) \) of \( W_k \):

\[
\Gamma \longmapsto \int_{\Gamma} e^{-W_k} \varphi_i.
\]

The following lemma shows that the set \( \{ e^{-W_k} \varphi_i \mid |i| \in \mathbb{Z}, |i| \leq N_k - p \} \) of relative cohomology classes forms a basis of \( \mathcal{F}_p H(W, \mu_d)_k \). It also shows that \( [\varphi_i] \in \Omega(W_k)^{\mu_d} \) corresponds to an element of the form \( \left( \varphi_i + \sum_{|j| < |i|} a_j \varphi_j \right) e^{-W_k} \) under (95).
Lemma A.2. — Under the isomorphism $H^N_k((\mathbb{C}^N)_k, W_k^{+\infty}) \cong H^{N_k-1}(W_k^{-1}(1))$, the class represented by $e^{-W_k}\varphi_i$ corresponds to the class $\eta_i$ in (94).

Proof. — Let $\Gamma$ be a Lefschetz thimble of $W_k$ in $H^N_{N_k}((\mathbb{C}^N)_k, W_k^{+\infty})$ and $C \in H^N_{N_k-1}(W_k^{-1}(1))$ be the corresponding cycle. (Note that $H^N_{N_k}((\mathbb{C}^N)_k, W_k^{+\infty}) \cong H^N_{N_k-1}(W_k^{-1}(1))$.) The image of $\Gamma$ under $W_k$ is assumed to be the positive real line. Then we have

$$\int_{\Gamma} e^{-W_k}\varphi_i = \int_0^{\infty} e^{-t}P(t)dt.$$  (96)

Here we set

$$P(t) := \int_{\Gamma \cap \{W_k(x) = t\}} \frac{\varphi_i}{dW_k} = \frac{1}{2\pi i} \int_{T} \frac{\varphi_i}{W_k(x) - t}$$  (97)

where $T$ is a circle bundle over $\Gamma \cap \{W_k(x) = t\}$. Using the homogeneity, one can deduce from the coordinate change $x_i \mapsto \frac{w_i}{d} x_i$ that

$$P(t) = t^{\vert i \vert - 1}P(1).$$

Therefore by (96),

$$\int_{\Gamma} e^{-W_k}\varphi_i = \Gamma(\vert i \vert)P(1) = \Gamma(1 - \{\vert i \vert\})P(\{\vert i \vert\} - 1)(1).$$  (98)

By differentiating (97) and setting $t = 1$, we find

$$\Gamma(1 - \{\vert i \vert\})P(\{\vert i \vert\} - 1)(1) = \int_C \eta_i.$$  (99)

The lemma follows from (98) and (99). □

Consider the tame deformation $W_{k,s}$ of $W_k$:

$$W_{k,s}(x) = W_k(x) + \sum_{i \in F_k} s_i x_i,$$

where $F_k := \{1 \leq j \leq N \mid \xi^{kw_j} = 1\}$ is the index set of co-ordinates on $(\mathbb{C}^N)_k$. For generic values of $s$, $W_{k,s}$ has only non-degenerate critical points (i.e. it is a Morse function). Let $z \in \mathbb{C}^\times$ and $\langle \cdot, \cdot \rangle : H^N_{N_k}((\mathbb{C}^N)_k, (W_{k,s}/z)^{+\infty}) \times H^N_{N_k}((\mathbb{C}^N)_k, (W_{k,s}/z)^{-\infty}) \rightarrow \mathbb{C}$ denote the intersection pairing (cf. (2)). Set

$$G_{ij}(s, z) := \{[e^{-W_{k,s}/z}\varphi_i], [e^{W_{k,s}/z}\varphi_j]\}.$$

This is a presentation of K. Saito’s higher residue pairing [61] by Pham [55]. The invariance of the pairing under the co-ordinate change $x_j \mapsto \lambda^{w_j/d} x_j$ shows the following.
Lemma A.3. — With respect to the degree \( \deg s := 1 - (w_i/d) \) and \( \deg z := 1 \), the function \( G_{\gamma}(s, z) \) is homogeneous of degree \(|i| + |j|\).

Lemma A.4. — The function \( G_{\gamma}(s, z) \) is regular at \( z = 0 \). Moreover

\[
G_{\gamma}(s, z) = (-1)^{\frac{N_i(N_j - 1)}{2}}(2\pi i z)^{N_i}(\text{Res}_{W_{k,i}}([\varphi_k], [\varphi_j]) + O(z)).
\]

Proof. — This is remarked in [55, 2ème Partie, Section 4.3, Remarque], but we include a proof for the convenience of the reader. Suppose that \( s \) is generic so that \( x \mapsto \Re(W_{k,i}(x)/z) \) is a Morse function. Let \( \Gamma_1^+, \ldots, \Gamma_L^+ \) (resp. \( \Gamma_1^-, \ldots, \Gamma_L^- \)) denote the Lefschetz thimbles emanating from the critical points \( \sigma_1, \ldots, \sigma_L \) of \( \Re(W_{k,i}/z) \) given by the upward (resp. downward) gradient flow. Choose an orientation of \( \Gamma_{i}^{\pm} \) such that \( \Gamma_{a}^{+} \cdot \Gamma_{b}^{-} = \delta_{ab} \). We have

\[
G_{\gamma}(s, z) = \sum_{a=1}^{L} \left( \int_{\Gamma_{a}^{+}} e^{-W_{k,i}/z_{\varphi_i}} \right) \cdot \left( \int_{\Gamma_{a}^{-}} e^{W_{k,i}/z_{\varphi_j}} \right).
\]

For a fixed argument of \( z \), we have the stationary phase expansion as \( z \to 0 \).

\[
\int_{\Gamma_{a}^{+}} e^{-W_{k,i}/z_{\varphi_i}} \sim \pm \frac{(2\pi z)^{N_i/2}}{\sqrt{\text{Hess} W_{k,i}(\sigma_a)}}(f_{i}(\sigma_a) + O(z)).
\]

Here we set \( \varphi_i = f_i(x) \wedge_{j \in \mathbb{F}_k} dx_j \), \( \text{Hess} W_{k,i}(\sigma_a) := \det((\partial_{x_i} \partial_{x_j} W_{k,i}), i \in \mathbb{F}_k) \) is the Hessian of \( W_{k,i} \) at \( \sigma_a \) and \( \pm \) is the sign depending on the orientation of \( \Gamma_{a}^{+} \). Therefore

\[
G_{\gamma}(s, z) \sim (-1)^{\frac{N_i(N_j - 1)}{2}}(2\pi i z)^{N_i}\left( \sum_{a=1}^{N} \frac{f_{i}(\sigma_a)f_{j}(\sigma_a)}{\text{Hess} W_{k,i}(\sigma_a)} + O(z) \right)
\]

where the lowest order term in the right-hand side equals the Grothendieck residue. The sign \((-1)^{\frac{N_i(N_j - 1)}{2}}\) comes from a local computation\(^{16}\) of the orientation. Since this holds for an arbitrarily fixed argument of \( z \), and \( G_{\gamma}(s, z) \) is holomorphic in \( z \in \mathbb{C}^{x} \), the conclusion follows for a generic \( s \). By analytic continuation, the same holds for all \( s \). \( \square \)

By Lemma A.3 and Lemma A.4, we have

\[
G_{\gamma}(0, z) = \begin{cases} 
0 & \text{if } |i| + |j| < N_k \\
(-1)^{\frac{N_i(N_j - 1)}{2}}(2\pi i z)^{N_i}\text{Res}_{W_{k}}([\varphi_i], [\varphi_j]) & \text{if } |i| + |j| = N_k
\end{cases}
\]

\(^{16}\) This comes from the local co-ordinate system centered at a critical point.
This shows the Hodge-Riemann bilinear relation:
\[(101) \quad (\mathcal{F}^p H(W, \mu_d)_k, \mathcal{F}^q H(W, \mu_d)_{d-k}) = 0 \quad \text{if } p + q > N_k.\]

For \(i, j\) such that \(|i|, |j| \in \mathbb{Z}\), we take lifts
\[ [e^{-W_i} \hat{\phi}_i] \in \mathcal{F}^p \cap \mathcal{F}^{N_k - p}, \quad [e^{-W_i} \hat{\phi}_j] \in \mathcal{F}^q \cap \mathcal{F}^{N_k - q} \]
which correspond to \([\varphi_i], [\varphi_j] \in \Omega W_k^\mu_d\) under the isomorphism \((95)\). When \(p + q > N_k\), the pairing \(\langle [e^{-W_i} \hat{\phi}_i], [e^{-W_i} \hat{\phi}_j]\rangle = 0\) vanishes by \((101)\). When \(p + q < N_k\), the pairing again vanishes because of the Hodge-Riemann bilinear relation \((101)\) for \(\mathcal{F}\). When \(p + q = N_k\), we have
\[ \langle [e^{-W_i} \hat{\phi}_i], [e^{-W_j} \hat{\phi}_j]\rangle = \frac{1}{d} \langle [e^{-W_i} \varphi_i], (1 - 1)\rangle [e^{W_j} \varphi_j] \quad \text{by } (4) \]
\[ = (-1)^{\frac{N_k(N_k-1)}{2}} (2\pi i)^{N_k} \frac{1}{d} \text{Res}_{W_i}([\varphi_i], (1 - 1)\rangle) \quad \text{by } (100). \]

The factor \((-1)^{\frac{k}{d}}\) comes from the map \(I^*\) in \((3)\). The proof of Proposition 2.1 is complete.

**Remark A.5.** — In the proof we observed that \(H(W, \mu_d)_k\) has a Hodge structure of weight \(N_k\). In order to make the weight compatible with the FJRW grading, we consider the Tate twist by \(\sum_{i=1}^{N} (k_i) - 1\) so that \(H(W, \mu_d)_k\) is of weight \(N_k + 2(\sum_{i=1}^{N} (k_i) - 1)\).

### Appendix B: Compatibility with FJRW setup

In [26] a factor \(f_g = |G|^e / \deg(st)\) multiplies all genus-\(g\) \(n\)-pointed invariants as well as all the homomorphisms
\[ \Lambda^{W,G}_{g,n} : H(W, G)^{\otimes n} \to H^*(\overline{M}_{g,n}; \mathbb{Q}), \]
\[ (\alpha_1, \ldots, \alpha_n) \mapsto f_g \text{st}_a \left( \overline{W}_{g,n,G}(W, (k_1, \ldots, k_n)) \right)^{\text{vir}} \cap \prod_{i=1}^{n} \alpha_i \]
defining a cohomological field theory via Poincaré duality. These operators embody all the relevant invariants via the definition
\[ \langle \tau_{\psi_{\alpha_1}}(\alpha_1), \ldots, \tau_{\psi_{\alpha_n}}(\alpha_n) \rangle_{W,G}^{g,n} = \int_{\overline{M}_{g,n}} \Lambda^{W,G}_{g,n}(\alpha_1, \ldots, \alpha_n) \prod_{i=1}^{n} \psi_{\alpha_i}^{k_i}. \]
The cycle $[\mathcal{W}_{g,n,G}(W, (k_1, \ldots, k_n))]^{\text{vir}} \cap \prod_{i=1}^n \alpha_i$ is a cycle in the moduli space of $(W, G)$-curves; in this paper $G$ equals $\mu_d$ and, in genus zero and for narrow state space entries, we may regard this as the top Chern class of the obstruction bundle. The degree of $\text{st}$ is simply the degree of the map forgetting the $(W, G)$-structure and retaining only the underlying coarse stable curve; for $(W, G) = (W, \mu_d)$ the morphism $\text{st}$ is the natural forgetful map $\text{Spin}^d_{g,n}(k_1, \ldots, k_n) \to \overline{\mathcal{M}}_{g,n}$. We have $\deg(\text{st}) = |G|^2 g - 1$ in general; therefore, $f_g$ equals $1/|G|^{g-1}$, and the setup of [26] is consistent with that of Witten’s original tentative treatment [70] of quantum singularity theory. In genus zero and for $G = \mu_d$, this amounts to an overall factor $d$ appearing also in [14, (14)].

We point out that all these different factors $f_g$ can be removed once we take into account that the pairing used (4) comes from orbifold Chen-Ruan cohomology (in its relative version) and acquires an overall factor $1/|G|$ equal to the degree of $BG$ over $\text{Spec } C$ (we recall that the pairing of [26] maps the pair $(\phi_k, \phi_i)$ to $\delta_{d-1,k+l}$ without any factor). In particular, removing the factor $f_0 = d$ in the definition of the genus-zero invariants does not change the quantum product: in the definition (17) of $T_i \circ T_j$, the factor $f_0$ is absorbed into $g_{k,l} = d \delta_{d-2,k+l}$. Furthermore, removing the factors $f_g$ from the cohomological field theory homomorphisms $\Lambda^{W,G}_{g,n}$ does not affect the composition axioms [26, (62), (64)].

Let $g = g_1 + g_2$; let $n = n_1 + n_2$; and let $\rho_{\text{tree}} : \overline{\mathcal{M}}_{g_1,n_1+1} \times \overline{\mathcal{M}}_{g_2,n_2+1} \to \overline{\mathcal{M}}_{g,n}$ be the gluing morphism. Then the forms

$$\widetilde{\Lambda}^{W,G}_{g,n}(\alpha_1, \ldots, \alpha_n) = \text{st}_* \left( [\mathcal{W}_{g,n,G}(W, (k_1, \ldots, k_n))]^{\text{vir}} \cap \prod_{i=1}^n \alpha_i \right) = \frac{\Lambda^{W,G}_{g,n}}{f_g}$$

satisfy the composition property stated in [26, (62)]

$$\rho_{\text{tree}}^* \widetilde{\Lambda}^{W,G}_{g,n}(\alpha_1, \alpha_2, \ldots, \alpha_n) = \sum_{\mu, \nu} g^{\mu,\nu} \Lambda^{W,G}_{g_1,n_1+1}(\alpha_1, \ldots, \alpha_{i_1}, \mu) \otimes \Lambda^{W,G}_{g_2,n_2+1}(\alpha_{i_1+1}, \ldots, \alpha_{i_1+n_2}, \nu),$$

for all $\alpha_i \in \text{H}(W, G)$, for $\mu$ and $\nu$ running through a basis of $\text{H}(W, G)$, and for $g^{\mu,\nu}$ denoting the inverse of the pairing $(,)$ with respect to the chosen basis. This happens because, by rescaling the pairing, we have multiplied $g^{\mu,\nu}$ by $|G|$; the canceled factors on the two sides of the above identity match

$$f_g = \frac{1}{|G|^{g-1}} = \frac{1}{|G|} \frac{1}{|G|^{g_1-1}} \frac{1}{|G|^{g_2-1}} = \frac{1}{|G|} f_{g_1} f_{g_2},$$

The same happens for the gluing morphism $\rho_{\text{loop}} : \overline{\mathcal{M}}_{g-1,n+2} \to \overline{\mathcal{M}}_{g,n}$. We have

$$\rho_{\text{loop}}^* \widetilde{\Lambda}^{W,G}_{g,n}(\alpha_1, \alpha_2, \ldots, \alpha_n) = \sum_{\mu, \nu} g^{\mu,\nu} \Lambda^{W,G}_{g-1,n+2}(\alpha_1, \ldots, \alpha_n, \mu, \nu).$$
Taking again into account that, in this paper, the matrix \((g^{\mu,\nu})\) has been multiplied by an overall factor \(|G|\), the cancellation of factors from the analogue identity \([26, (64)]\) yields the same quantity on both sides

\[
f_g = \frac{1}{|G|} \frac{1}{|G|^{\frac{2}{2}}} = \frac{1}{|G|} f_g^{-1}.\]

REFERENCES

1. D. Abramovich, T. Graber, and A. Vistoli, Gromov-Witten theory of Deligne-Mumford stacks, *Am. J. Math.*, 130 (2008), 1337–1398. arXiv:math/0603151.
2. D. Abramovich and A. Vistoli, Compactifying the space of stable maps, *J. Am. Math. Soc.*, 15 (2002), 27–75 (electronic).
3. P. S. Aspinwall, D-branes on Calabi-Yau manifolds, in *Progress in String Theory*, pp. 1–152, World Scientific Publishing, Hackensack, 2003.
4. V. V. Batyrev, Dual polyhedra and mirror symmetry for Calabi-Yau hypersurfaces in toric varieties, *J. Algebr. Geom.*, 3 (1994), 493–535.
5. L. A. Borisov and R. Paul Horja, Mellin-Barnes integrals as Fourier-Mukai transforms, *Adv. Math.*, 207 (2006), 876–927.
6. L. A. Borisov and R. Paul Horja, On the better-behaved version of the GKZ hypergeometric system. arXiv:1011.5720.
7. R.-O. Buchweitz, Maximal Cohen-Macaulay modules and Tate cohomology over Gorenstein rings, 1986. https://tspace.library.utoronto.ca/bitstream/1807/16682/1/maximal_cohen-macaulay_modules_1986.pdf.
8. R.-O. Buchweitz, G.-M. Greuel, and F.-O. Schreyer, Cohen-Macaulay modules on hypersurface singularities. II, *Invent. Math.*, 88 (1987), 165–182.
9. A. Canonaco and R. L. Karp, Derived autoequivalences and a weighted Beilinson resolution, *J. Geom. Phys.*, 58 (2008), 743–760.
10. W. Chen, Y. Ruan, Orbifold Gromov-Witten theory, in *Orbifolds in Mathematics and Physics* (Madison, WI, 2001), Contemp. Math., vol. 310, pp. 25–85, Am. Math. Soc., Providence, 2002.
11. A. Chiodo, Stable twisted curves and their \(r\)-spin structures (Courbes champêtres stables et leurs structures \(r\)-spin), *Ann. Inst. Fourier (Genoble)*, 58 (2008), 1635–1689. arXiv:math.AG/0603687.
12. A. Chiodo and J. Nagel, in preparation.
13. A. Chiodo and Y. Ruan, Landau-Ginzburg/Calabi-Yau correspondence for quintic three-folds via symplectic transformations, *Invent. Math.*, 182 (2010), 117–165. arXiv:0908.0908.
14. A. Chiodo and Y. Ruan, LG/CY correspondence: the state space isomorphism, *Adv. Math.*, 227 (2011), 2157–2188. arXiv:0908.0908.
15. A. Chiodo and Y. Ruan, A global mirror symmetry framework for the Landau-Ginzburg/Calabi-Yau correspondence, *Annales de l’Institut Fourier*, to appear. http://www-fourier.ujf-grenoble.fr/~chiodo/framework.
16. A. Chiodo and D. Zvonkine, Twisted \(r\)-spin potential and Givental’s quantization, *Adv. Theor. Math. Phys.*, 13 (2009), 1335–1369. arXiv:0711.0339.
17. T. Coates, A. Corti, H. Iritani, and H.-H. Tseng, Computing genus-zero twisted Gromov-Witten invariants, *Duke Math. J.*, 147 (2009), 377–438. arXiv:math.AG/0702234.
18. T. Coates, A. Corti, H. Iritani, and H.-H. Tseng, in preparation.
19. T. Coates, A. Corti, Y.-P. Lee, and H.-H. Tseng, The quantum orbifold cohomology of weighted projective spaces, *Acta Math.*, 202 (2009), 139–193.
20. T. Coates, A. Gholampour, H. Iritani, Y. Jiang, P. Johnson, and C. Manolache, The quantum Lefschetz hyperplane principle can fail for positive orbifold hypersurfaces, *Math. Res. Lett.*, 19 (2012), 997–1005. arXiv:1202.2754.
21. T. Coates and A. Givental, Quantum Riemann-Roch, Lefschetz and Serre, *Ann. Math.* (2), 165 (2007), 15–53.
22. T. Dyckerhoff, Compact generators in categories of matrix factorizations. *Duke Math. J.*, 159 (2011), 223–274. arXiv:0904.3413.
56. A. Polishchuk and A. Vaintrob, Chern characters and Hirzebruch-Riemann-Roch formula for matrix factorizations. 
*Duke Math. J.*, **161** (2012), 1863–1926. arXiv:1002.2116.

57. A. Polishchuk and A. Vaintrob, Matrix Factorizations and Cohomological Field Theory. arXiv:1105.2903.

58. A. Pressley and G. Segal, *Loop Groups*, Oxford Mathematical Monographs, Clarendon Press/Oxford University Press, New York, 1986. Oxford Science Publications.

59. T. Reichelt, A construction of Frobenius manifolds with logarithmic poles and applications, *Commun. Math. Phys.*, **287** (2009), 1145–1187.

60. M. A. Rose, A reconstruction theorem for genus zero Gromov-Witten invariants of stacks, *Am. J. Math.*, **130** (2008), 1427–1443.

61. K. Saito, The higher residue pairings $K^{-1}_{\mathbb{F}}$ for a family of hypersurface singular points, in *Singularities, Part 2 (Arcata, Calif., 1981)*, Proc. Sympos. Pure Math., vol. 40, pp. 441–463, Am. Math. Soc., Providence, 1983.

62. E. Segal, Equivalence between GIT quotients of Landau-Ginzburg B-models, *Commun. Math. Phys.*, **304** (2011), 411–432.

63. P. Seidel and R. Thomas, Braid group actions on derived categories of coherent sheaves, *Duke Math. J.*, **108** (2001), 37–108.

64. J. Steenbrink, Intersection form for quasi-homogeneous singularities, *Compos. Math.*, **34** (1977), 211–223.

65. B. Toen, Théorèmes de Riemann-Roch pour les champs de Deligne-Mumford, *K-Theory*, **18** (1999), 33–76.

66. H.-H. Tseng, Orbifold quantum Riemann-Roch, Lefschetz and Serre, *Geom. Topol.*, **14** (2010), 1–81.

67. C. Vafa and N. P. Warner, Catastrophes and the classification of conformal theories, *Phys. Lett. B*, **218** (1989), 51–58.

68. J. Walcher, Stability of Landau-Ginzburg branes, *J. Math. Phys.*, **46** (2005), 082305, 29.

69. E. Witten, Phases of $\mathcal{N}=2$ theories in two dimensions, *Nucl. Phys. B*, **461** (1994), 159–222.

70. E. Witten, *Algebraic Geometry Associated with Matrix Models of Two-Dimensional Gravity*, Topological Models in Modern Mathematics, Publish or Perish, Stony Brook, New York, 1991.

A. C.
Institut de Mathématiques de Jussieu, UMR 7586 CNRS,
Université Pierre et Marie Curie,
Case 247, 4 Place Jussieu,
75252 Paris cedex 05, France
chiodoa@math.jussieu.fr

H. I.
Department of Mathematics, Graduate School of Science,
Kyoto University,
Kitashirakawa-Oiwake-cho, Sakyoku,
Kyoto, 606-8502, Japan
iritani@math.kyoto-u.ac.jp

Y. R.
Department of Mathematics,
University of Michigan,
Ann Arbor, MI 48109-1109, USA
ruan@umich.edu

Manuscrit reçu le 12 janvier 2012
Manuscrit accepté le 10 mai 2013
publié en ligne le 20 juin 2013.