The growth and assembly of a massive galaxy at $z \sim 2$

N. A. Hatch$^{1,*}$, R. A. Overzier$^{2}$, J. D. Kurk$^{3}$, G. K. Miley$^{1}$, H. J. A. Röttgering$^{1}$ and A. W. Zirm$^{4}$

$^{1}$Leiden Observatory, University of Leiden, P.B. 9513, Leiden 2300 RA, The Netherlands
$^{2}$Max-Planck-Institut für Astrophysik, Karl-Schwarzschild-Str. 1, D-85748, Garching, Germany
$^{3}$Max-Planck-Institute für Astronomie, Königstuhl 17, 69117 Heidelberg, Germany
$^{4}$Department of Physics and Astronomy, Johns Hopkins University, 3400 North Charles Street, Baltimore, MD21218

ABSTRACT

We study the stellar mass assembly of the Spiderweb Galaxy (MRC 1138-262), a massive $z = 2.2$ radio galaxy in a protocluster and the probable progenitor of a brightest cluster galaxy. Nearby protocluster galaxies are identified and their properties are determined by fitting stellar population models to their rest-frame ultraviolet to optical spectral energy distributions. We find that within 150 kpc of the radio galaxy the stellar mass is centrally concentrated in the radio galaxy, yet most of the dust-uncorrected, instantaneous star formation occurs in the surrounding low-mass satellite galaxies. We predict that most of the galaxies within 150 kpc of the radio galaxy will merge with the central radio galaxy by $z = 0$, increasing its stellar mass by up to a factor of $\sim 2$. However, it will take several hundred Myr for the first mergers to occur, by which time the large star formation rates are likely to have exhausted the gas reservoirs in the satellite galaxies. The tidal radii of the satellite galaxies are small, suggesting that stars and gas are being stripped and deposited at distances of tens of kpc from the central radio galaxy. These stripped stars may become intracluster stars or form an extended stellar halo around the radio galaxy, such as those observed around cD galaxies in cluster cores.

Key words: galaxies: elliptical and lenticular, cD – galaxies: individual: MRC 1138-262 – galaxies: evolution – galaxies: high-redshift

1 INTRODUCTION

High redshift radio galaxies (HzRGs) are among the most luminous and massive known galaxies (Seymour et al. 2007 and references therein). They contain vast reservoirs of atomic and molecular gas, and are commonly embedded in giant ionized gas halos. A number of well studied HzRGs have been found to be undergoing violent activity such as extreme star formation and merging. They contain active galactic nuclei (AGN) implying the co-evolution of their supermassive black holes and stellar hosts. Some radio galaxies have been found to be located in dense regions of the early Universe that may collapse to form clusters of galaxies with masses greater than $10^{14} M_\odot$ when extrapolated to $z = 0$ (Venemans et al. 2007). Therefore HzRGs are likely to be the progenitors of low-redshift brightest cluster galaxies (BCGs; Milew & De Breuck 2008; Zirm et al. 2005, Pentericci et al. 1997).

With a mass of $\sim 10^{12} M_\odot$, the radio galaxy MRC 1138-262 at $z = 2.2$ (Fig. 1; Pentericci et al. 1997, Milew et al. 2004, Overzier et al. 2007), also known as the ‘Spiderweb Galaxy’, is one of the most massive galaxies at $z \sim 2$ (Seymour et al. 2007). It lies in a protocluster identified by an overdensity of Ly$\alpha$ and H$\alpha$ emitters, extremely red objects and sub-mm bright galaxies (Pentericci et al. 2000, Kurk et al. 2004, Kodama et al. 2007, Stevens et al. 2003, Zirm et al. 2008). The radio galaxy is 2 mag brighter in $K_s$ than any other protocluster galaxy candidate within a 1.5 Mpc physical radius (Kurk et al. 2004). It is therefore highly likely that this radio galaxy is the central galaxy within a massive dark matter halo at the centre of the protocluster, and that the nearby galaxies may become satellite galaxies, or eventually merge with it to form a brightest cluster galaxy.

Hubble Space Telescope (HST) images of this radio source in Miley et al. (2004) show that MRC 1138-262 is surrounded by tens of UV-bright clumps and galaxies. These observations agree qualitatively with predictions of hierarchical galaxy formation theory. In this work we examine the properties of the nearby galaxies in more detail, in order to quantify the future growth and mass assembly of the radio galaxy through the accretion of these galaxies. We combine high resolution rest-frame UV and optical data to identify protocluster galaxies that are likely to merge with the radio galaxy. By comparing their photometry to stellar population models, we estimate the stellar mass of these galaxies. We then use the observed stellar mass distribution as initial conditions to examine the future growth and assembly of the radio galaxy.

In section 2 we describe the observations and data reduction
and section 3 describes how the protocluster galaxies are identified. Section 4 gives details of the spectral energy distribution (SED) fitting, through which we derive the masses of the galaxies. These properties are discussed and compared to semi-analytic models in section 5. In section 5 we use simple analytic approximations of the merging timescale to examine the evolution of MRC 1138-262. Throughout this work we use $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.27$, and $\Omega_{\Lambda} = 0.73$ (Spergel et al. 2003), and at the redshift of MRC 1138-262 ($z=2.156$) the linear scale is $8.4 \text{kpc arcsec}^{-1}$. All magnitudes are AB magnitudes.

2 OBSERVATIONS

The protocluster surrounding MRC 1138-262 was observed with the Advanced camera for Surveys (ACS) instrument on the HST through the F475W ($g_{475}$) filter for 9 orbits, and 10 orbits through the F814W ($i_{814}$) filter. The data were reduced using the ACS pipeline (APSIM; Blakeslee et al. 2003). Details of the data reduction are given in Miley et al. (2000) and Hatch et al. (2008). The protocluster was also observed with NICMOS camera 3 on HST for 3 orbits through each of the F110W ($J_{110}$) and F160W ($H_{160}$) filters. Further discussion of the NICMOS observations can be found in Zirm et al. (2008).

The HST observations were supplemented with a 3 hour exposure in $U_n$-band from the LRIS-B instrument on the Keck telescope (P. W. van Breugel), and near-infrared (near-IR) images from the Very Large telescope (VLT) ISAAC instrument: $Ks$–band (1.6 hour) and a narrow-band image centred on 2.07 $\mu$m (4.8 hour). The $U_n$–band data has a full width at half maximum (FWHM) of 1.2 arcsec, and both near-IR images have a FWHM of 0.45 arcsec, measured from bright, unsaturated stars that lie close to the radio galaxy. Details of the reduction of the near-IR observations can be found in Kurk et al. (2003). H$\alpha$ emission from the protocluster at $z=2.156$ lies within the bandwidth of the narrow-band NB207 filter. Both an H$\alpha$ emission-line image of the protocluster and an emission-line-free $Ks$ continuum image were created from the near-IR dataset.

The images were registered and matched in pixel-scale (0.2 arcsec pixel$^{-1}$) and resolution to the $Ks$ image. The $3\sigma$ image depths of the images are 27.7 mag in $g_{475}$, 27.7 mag in $i_{814}$, 26.2 mag in $J_{110}$, 26.6 mag in $H_{160}$, 24.4 mag in $Ks$ and 26.8 mag in $U_n$. Magnitudes were measured using SExtractor (Bertin & Arnouts 1996) in double-image mode, using a "white-light" image as a detection image. The detection image was created by combining the high signal-to-noise $g_{475}$, $i_{814}$, $J_{110}$ and $H_{160}$ HST images. Regions with 7 or more connecting pixels which had a signal-to-noise ratio greater than 4 times the local background noise were identified and isophotal magnitude were measured.

The point spread function (PSF) of the $U_n$–band data is more than twice as large as the PSF of the images. To preserve the high resolution of the HST and $Ks$ images we scale the $U_n$ magnitudes by a factor to account for the light dispersed beyond the galaxy isophotal detection area by the larger PSF. The $g_{475}$ was smoothed with a Gaussian of FWHM 1.2 arcsec to match the resolution of the $U_n$ image. Magnitudes were obtained from both the $g_{475}$ and $U_n$ images using the original detection image and multiplied by the ratio of flux within the isophotal area of the 0.45 arcsec resolution $g_{475}$ image to the 1.2 arcsec resolution $g_{475}$ image. Objects which are fainter in the 0.45 arcsec resolution image than the 1.2 arcsec resolution $g_{475}$ image are blended and no $U_n$ magnitude was assigned to these objects. All magnitudes were corrected for Galactic extinction of $A_g = 0.172$ (Schlegel et al. 1998).

2.1 Emission-line contribution

Ly$\alpha$ emitted from galaxies in the protocluster lies within the bandwidth of the $U_n$ filter, although the filter transmission at this wavelength is only 30% of the maximum filter transmission. Kurk et al. (2004) lists the bright Ly$\alpha$ emitters close to the radio galaxy. Galaxy 10 (see Fig. 4) has an equivalent width (EW) greater than 1000$\AA$. The $U_n$ magnitude for this galaxy is likely to be dominated by the Ly$\alpha$ emission. We therefore assume the observed magnitude is an upper limit on the $U_n$ continuum emission. The other Ly$\alpha$ emitters have EWS of 30-100$\AA$, therefore a maximum of 5% of the emission in the $U_n$–band results from Ly$\alpha$ line emission and we do not include any correction for this.

Both the CIV and He II emission lines fall within the $g_{475}$ bandwidth, and contribute greatly to the diffuse intergalactic light (Hatch et al. 2008), but spectroscopy shows that neither line is detected in the majority of the Ly$\alpha$ emitting galaxies close to MRC 1138-262 (Kurk 2003). Only the radio galaxy and galaxy 10 have large combined CIV and He II EWS of 101$\AA$ and 87$\AA$ respectively and the $g_{475}$ magnitudes are reduced by 22.5% and 20% respectively.

H$\beta$ and [OII] emitted from the protocluster galaxies falls within the $H_{160}$ bandwidth. The contribution in $H_{160}$ from these lines is estimated using the H$\alpha$ image and observations of MRC 1138-262 from the integral field unit spectrometer SINFONI, which show that H$\alpha$+[NII] $> 2.4 \times$ [OII]+H$\beta$ (N. Nesvadba – private communication). The line emission from the galaxies is likely to be fainter than the line emission from the halo, therefore a maximum of 5% of the $H_{160}$ flux can be due to line emission. As we do not know exactly how much line emission should be removed from each galaxy, we do not remove this component. The H$\alpha$ emission was removed from the $Ks$ image using the narrow-band NB207 image.

2.2 Removing the nuclear contribution

The nuclear radio emission is associated with galaxy 1 and we refer to this galaxy as the radio galaxy (see Fig. 2; Penetier et al. 1997). To remove the point source contribution from the active galactic nucleus, all galaxies, except for the radio galaxy, are masked out from the original high-resolution HST images, and the resulting image is fitted with an ellipsoid model to obtain the best-fitting centroid, position angle and axial ratio. The azimuthally averaged radial profile is fitted with a classical de Vaucouleurs profile plus a point source smoothed to the resolution of the image. The centroid of the point source is not constrained to lie at the centre of the large-scale ellipsoid. The point source contribution to the $Ks$, $H_{160}$, $J_{110}$ and $i_{814}$ images is 15%, 30%, 23%, 9% respectively and the galaxy profile in $g_{475}$ and $U_n$–bands is consistent with no point source contribution. Most AGN have nuclear starbursts and therefore a small percentage of star formation is likely being removed together with the point source.

The radio galaxy radial profile in $J_{110}$ and $H_{160}$ is well described by a central point source plus a de Vaucouleurs profile with a characteristic radius of 0.6 arcsec (5 kpc – comparable to low redshift cD and giant Elliptical galaxies). However, in the rest-frame UV light ($i_{814}$ and $g_{475}$) the radio galaxy is not well described by the de Vaucouleurs plus point source profile at large radii because
of the large contribution of extended diffuse light (see Hatch et al. 2008).

3 SELECTING PROTOCLUSTER GALAXIES

We wish to study the central region of the protocluster, the radio galaxy and its neighbouring satellite galaxies within 150 kpc. Total extent of the image is 300 kpc on each side. Galaxies identified to be within the protocluster are marked and labelled in Fig. Note the bridge of red light connecting galaxies 17 and 18 possibly arising from tidal stripping.

Eight of the 64 objects are confirmed to be protocluster galaxies through the detection of Lyα emission (galaxies 1, 5, 6, 7, 9, 10, 11, 15: Kurk 2003; Kurk et al. 2004; see Fig. 2). Five objects are observed to emit Hα emission, galaxies 1, 10, 11, 12 are identified by Kurk et al. (2004) and we identify a further galaxy (ID 8) at 11:40:48.37 – 26:29:05.2 (J2000) with an EW of 153Å. In total 9 satellites of the radio galaxy are identified by line emission detections.

Selecting protocluster galaxies through line emission alone biases the sample toward star forming galaxies. At the redshift of the protocluster the Balmer and 4000 Å break falls within the J110 filter, therefore we can make a simple colour cut to select candidate red protocluster galaxies. Zirm et al. (2008) show that there is a significant overdensity of galaxies in this protocluster field with colours 1.1 < J110 − H160 < 2.1. Using the same colour cut we identify 7 red galaxies (galaxies 1, 2, 4, 12, 13, 14 and 16) two of which are also Hα emitters.

We include 3 additional galaxies in our catalogue of satellite galaxies based on their proximity to the radio galaxy and their peculiar morphologies. Galaxies 17 and 18 have a bridge of red light connecting them which corresponds to a surface brightness enhancement in the Lyα halo (see Fig. 1 and 2) suggesting that they are an interacting system, but the photometric redshift estimation code BPZ (Benítez 2000) assigns redshifts of 3.6 and 0.3 respectively with low probability. Galaxy 19 has a peculiar linear morphology and lies close in projection to the radio galaxy. A concentration of chain and tadpole galaxies were found around the radio galaxy (Miley et al. 2006) suggesting that they are located in the protocluster core. We include galaxies 17, 18 and 19 in the subsequent analysis but note that these galaxies have not been formally identified with the protocluster.

Photometric redshifts of all objects were obtained using the bayesian photometric redshift estimation code BPZ (Benítez 2000), utilizing the CWWSB template set as described in Benítez (2000) and 2 additional simple stellar population models with ages of 5 and 25 Myr (Coe et al. 2004) from Bruzual & Charlot (2003). We used a prior based on the galaxy population of the Northern Hubble Deep Field and imposed an additional weak prior of a cluster at z = 2.2 with 1% of the galaxies in the field lying within the cluster. This is a weak additional prior since 10% of the galaxies in the 0.283 arcmin² field are identified as protocluster galaxies from line emission.

The Hα emitters labelled 8 and 12, and all protocluster galaxies selected by a red colour-cut have photometric redshifts within $\Delta z = 0.3$ of the redshift of the radio galaxy. One further galaxy (3) has a photometric redshift of 2.2 with Bayesian Odds $O = 0.9$. Including the cluster prior shifted the best fit redshifts of galaxies 13, 14 and 16 from $z \sim 1.4$ to $z = 2.2$, although $z = 2.2$ lay within the 1σ confidence intervals whether the cluster prior was included or not. Increasing the cluster prior such that 50% of the galaxies in the field lie within the cluster had no further effect and did not increase the number of galaxies with photometric redshifts within $\Delta z = 0.3$ of 2.2. The photometric redshifts of the
blue galaxies could not be constrained using BPZ. Galaxies with high equivalent width of Lyα generally have smaller UV colours across the Lyman break than galaxies with no Lyα emission. Also, although the selection based on U-G is a powerful technique of isolating ‘BX’ galaxies at $z \sim 2.2 \pm 0.3$ (see Reddy et al. 2005), our U band data does not reach deep enough to constrain the colour across the Lyman break accurately.

Within 150 kpc of the radio galaxy, we find a total of 19 galaxies in the protocluster. It should be noted that the number of galaxies seen in the $I_{814}$ images of Miley et al. (2006) and Hatch et al. (2008) is greater than presented in this work. The lower resolution of the data used in this work removes many of the smaller clumps of bright emission. The large number of the small clumps seen near the radio galaxy in the $I_{814}$ and $g_{475}$ images suggest they are located in the same region as the radio galaxy although no redshift can be determined. Thus the 19 detected galaxies should be considered the observable subset of a larger population of protocluster galaxies.

4 STELLAR POPULATION SYNTHESIS MODELING

A number of works have derived the stellar population properties of galaxies using stellar population synthesis techniques. We follow a procedure similar to that of Papovich et al. (2001), Shapley et al. (2005) and Kriek et al. (2008) to determine stellar population properties by fitting the broad-band photometry with stellar population models. Our aim was to determine the age, intrinsic star formation rate, $E(B-V)$, star formation history and stellar mass of each galaxy within 150 kpc of the radio galaxy.

We compare the spectral energy distribution (SED) of the galaxies with stellar population templates created using Bruzual & Charlot (2003) models. We used models with solar metallicity, and a Chabrier (2003) initial mass function (IMF) extending from 0.1 to 100 M$_\odot$. We use the Padova 1994 stellar population evolution tracks to form single exponentially declining star formation histories with characteristic e-folding times ($\tau$) of 1, 0.5, 0.1, 0.05, and 0.01 Gyr, and a model with a constant star-forming history. We consider models with 31 ages between $10^6$ yrs since the onset of star formation and 3.0 Gyr, the age of the Universe at $z = 2.156$, the redshift of MRC 1138-262; and 21 extinction values between $A_V = 0$ and 4 mag, assuming the Calzetti reddening law. We compare the photometry to the stellar population synthesis templates and determine the best-fit parameters by minimum-$\chi^2$ techniques. Confidence intervals due to the photometric errors are derived by constructing 500 synthetic realisations of each galaxy, varying the photometry by a random amount drawn from the normal photometric error distribution. The stellar population models are compared to the synthetic data sets to derive new best-fitting models. $1\sigma$ confidence intervals of the best fit parameters are derived from the minimum and maximum values enclosing 68 per cent of the simulations. We also derive best-fit parameters without the $U$-band information by comparing the $g_{475}$ through $K_s$-band photometry to the stellar population models. The results of all the alternative fits are within the errors given in Table 1.

Stellar templates created with single exponentially decreasing or constant star formation histories provide lower limits to the galaxy mass, particularly if the galaxy has a very young age and active star formation. Such galaxies have relatively small mass-to-light ratios as the observed light is dominated by bright young stars, which may hide an older stellar population, that could potentially be more massive than the brighter young population. To estimate the amount of mass that may exist in an older stellar population we fit the photometry with a two component star formation history. One component is fixed to be a 3 Gyr population with $\tau = 0.01$ Gyr, whilst the other model is fixed to be a 1 Myr old constant star forming population. The galaxy mass estimated from this two-component model is an upper limit on the true stellar mass. $1\sigma$ confidence intervals are drawn from 500 Monte Carlo simulations as described above.

The best-fit parameters depend systematically on the choice of metallicity and IMF, for example a Salpeter IMF gives mass-to-light ratios that are systematically higher by $\sim 1.5$ than the Chabrier IMF. Larger metallicities result in redder SEDs, and thus allowing the metallicity to decrease, can increase the best-fit extinction, and decrease the age of the stellar population. A full discussion of how the best-fit galaxy parameters depend on the IMF and the metallicity is given in [Papovich et al. 2001].

The parameters derived from SED fitting are subject to large systematic uncertainties and degeneracies and we refer to Shapley et al. (2005) for details. The best-fit age, intrinsic star formation rate and extinction strongly depend on the assumed star formation history. The data are unable to constrain the star formation histories across our range of parameter space, for all the selected galaxies. We therefore infer that the derived age, extinction and intrinsic star formation rate are unreliable, and will not be discussed further. In the following discussion we refer only to the derived stellar mass (columns 3 and 4 of Table I), and the dust-uncorrected star formation rate (column 2), determined using the observed $g_{475}$ luminosity and

$$SFR (M_{\odot} yr^{-1}) = \frac{L_{1500}}{1.5 \times 10^{28}} (\text{erg s}^{-1} \text{Hz}^{-1})$$

(Salim et al. 2007).

Stellar masses do not suffer the same large systematic uncertainties due to the unconstrained star formation history [Shapley et al. 2005], and therefore the derived masses are likely to be more robust. We note that Shapley et al. (2005) use near-infrared photometry, which means the derived masses are more reliable. The rest-frame $V$-band is longest wavelength in this work, and therefore we expect larger uncertainties. Although Spitzer mid-infrared images are available, the high density of the galaxies in this region, and the low resolution of the data, does not allow us to robustly measure rest-frame near-infrared magnitudes of the individual galaxies.

Recent work has focused on the contribution to the rest-frame near-infrared light of thermally pulsating asymptotic giant branch stars, which occurs 0.2–2 Gyr after the onset of star formation [Maraston et al. 2006; Bruzual & Charlot 2003]. The redder band we use in this work is equivalent to the rest-frame $V$, and the contribution of thermally pulsating asymptotic giant branch stars at this wavelength is not important. We have compared our results from the Bruzual & Charlot (2003) models to the new Charlot & Bruzual (2007) models which include the thermally pulsating asymptotic giant branch stars. We find no significant change in the results given in Table I in particular the stellar masses are within the 1$\sigma$ errors quoted.

We note that masses derived from multiband photometry are subject to systematic uncertainties from the assumed IMF [Conroy et al. 2008]. The uncertainty on the assumed IMF may drastically change the derived masses, possibly by more than a factor of 10. However, the relative masses of galaxies, which are of particular importance to this work, are derived with much greater certainty.
The growth and assembly of a massive galaxy

Figure 3. Broad-band photometry (black filled circles) and 1σ errors are shown for the radio galaxy (ID 1) and 18 satellite galaxies. The point source contamination from the AGN has been removed from the radio galaxy. The best-fit stellar population models are shown in grey, the diamond symbols are the convolution of the best-fit stellar population models with the filter transmission curves.

5 RESULTS

5.1 Properties of the galaxies

5.1.1 Stellar mass

The nucleus of the radio emission is associated with galaxy 1 (Pentericci et al. 1997). The total stellar mass of this galaxy obtained from a fit to a single stellar population model is $1.1 \pm 0.2 \times 10^{12} M_\odot$ which is consistent with the upper limit obtained from the rest-frame near-infrared light (Seymour et al. 2007). This galaxy dominates the central 150 kpc in terms of mass, comprising $85 \pm 5$ per cent of the total mass. If one assumes that the mass of the satellite galaxies is closer to the upper mass limit (listed in column 3 of Table[1] the fraction of mass in the radio galaxy reduced to $60 \pm 13$ per cent. Rest-frame UV and optical light may be scattered from the central active nucleus contaminating the stellar light of galaxy 1, resulting in an overestimate of the stellar mass. Polarization measurements are essential to determine the fraction of scattered quasar light within each band. The upper mass limit of galaxy 1 is an order of magnitude larger than previous measurements. We therefore always assume in the following discussion, that the mass of this radio galaxy is that of the lower mass limit ($1 \times 10^{12} M_\odot$) which is consistent with the measurement of Seymour et al. (2007).

Galaxy 1 is at least $\sim 10$ times more massive than any other galaxy within 150 kpc. The sample we have selected is unlikely to be a complete inventory of satellite galaxies. From a fit of the observed $K_s$ mag and the derived upper mass limit, we find that the 3σ $K_s$-band detection limit of 24.4 mag corresponds to a mass of $\sim 6 \times 10^9 M_\odot$. Our sample is unlikely to be missing any satellite galaxies more massive than this mass limit. We identify galaxy 1 as the progenitor of the BCG, which already dominates the mass budget of the protocluster core, and we refer to this galaxy as the radio galaxy or the central galaxy. The remaining protocluster galaxies within 150 kpc of the radio galaxy are referred to as the satellite galaxies.
This galaxy [Stevens et al. 2003], with a 850μm flux density of 6.7±2.4 mJy. Thus the total star formation rate of the region close to the radio galaxy is potentially much higher than given by the dust-uncorrected UV luminosity. The galaxies have red UV colours (see Fig. 5 in section 5.2.2 implying extinctions in the approximate range 0.1 < E(B−V)< 0.5 depending on the age, star formation history and metallicity of the stellar population. Correcting the UV luminosity for this amount of dust extinction would bring the star formation rates derived from the observed UV and sub-mm emission into agreement. The 850μm emission from this location is spatially elongated along a position angle of 72° [Stevens et al. 2003]. The satellite galaxies with extreme star formation and large amounts of extinction (galaxies 5, 6, 7, 9, 10, 11 and 15) are elongated approximately East-West along the sky stretching 15 arcsec from end-to-end. Therefore the extended sub-mm emission results from source confusion and is the sum of the combined star formation occurring in the radio and satellite galaxies. High spatial resolution sub-mm observations have deconvolved the extended sub-mm emission from the high redshift galaxy 4c40.7 into two separate components indicating merger induced star formation [Vilyon et al. 2008]. These optical and near-infrared observations show the star formation occurs in multiple compact objects and high resolution sub-mm may be able separate the sub-mm emission into the various galaxy components.

5.1.3 Spatial distribution of the satellite galaxies

Almost all of the satellite galaxies lie within the large Lyα halo that surrounds the radio galaxy, even those galaxies with little or no star formation. Since our observations show a single snapshot in time, most of the satellite galaxies' orbits must lie within the Lyα halo. This gives an approximate orbit radius of 100 kpc. The asymmetric Lyα halo extension to the East also contains satellite galaxies. This strong spatial correlation implies a connection between the Lyα halo and the satellite galaxies. The Lyα halo may originate from gas stripped from the satellites, or the galaxies may excite the gas, possibly by shocks from their movement through an existing gas reservoir, or UV emission from hot young stars. The total Lyα luminosity of the halo is 2.6×10⁸ erg s⁻¹ and would require a star formation rate of ~2500 M⊙ yr⁻¹ to ionize it. The UV star formation rates are more than an order of magnitude too small to ionize the halo, although the bright sub-mm emission suggests a large amount of obscured star formation. It is possible that the Lyα halo is ionized by the UV emission from young hot stars produced during this phase of extreme star formation, or is a mixture of both stellar and AGN photoionization [Villar-Martin et al. 1997; Villar-Martin et al. 2007].

5.2 Comparison of the Spiderweb system to simulations

We compare our results to two different semi analytic models (SAMs): De Lucia & Blaizot 2007 (DLB07) and Font et al. 2008 (F08) which are based on the dark matter only Millennium simulation1. We construct catalogues from the SAMs that most resemble our own selection criteria. We select all galaxies within 10 Mpc of the most massive (stellar mass) 10 galaxies at z = 2.2, then make a further cut selecting only those galaxies that lie within a projected distance of 150 kpc from the massive galaxy. We take an average of 3 orthogonal projections of each of the 10 galaxies.

Table 1. The mass is derived from fitting the photometry to a single exponentially declining star formation history. The mass upper limit is derived from a two-model fit to the photometry, in which one model is maximally old, i.e., 3 Gyrs and the other is maximally young (1 Myr). Column 5 lists the detection methods by which the galaxy was selected to be in the protocluster. Lyα, Hα are objects which have an excess of line-emission placing them at the same redshift as the radio galaxy, BB indicates galaxies with strong Balmer breaks inferred from large observed J₁₁₀−H₁₆₀ colours.

| ID | SFR (M⊙ yr⁻¹) | Mass (10⁸ M⊙) | Mass (upper limit) (10⁹ M⊙) | Detection method |
|----|---------------|---------------|-----------------------------|-----------------|
| 1  | 16.0 ± 0.7    | 1100 ± 200    | 9900 ± 32000                | Lyα, Hα, BB     |
| 2  | 0.0 ± 0.4     | 12 ± 0.3      | 16 ± 0.6                    | BB              |
| 3  | 2.2 ± 0.5     | 11 ± 8        | 105 ± 13                    | Photo-z         |
| 4  | 0.0 ± 0.4     | 31 ± 24       | 34 ± 3                      | BB              |
| 5  | 5.2 ± 0.5     | 7 ± 0.4       | 15 ± 3                      | Lyα             |
| 6  | 4.4 ± 0.5     | 9 ± 3         | 37 ± 15                     | Lyα             |
| 7  | 3.8 ± 0.5     | 3 ± 0.6       | 6 ± 0.9                     | Lyα             |
| 8  | 1.2 ± 0.5     | 5 ± 0.5       | 72 ± 6                      | Hα              |
| 9  | 7.2 ± 0.5     | 5 ± 2         | 9 ± 5                       | Lyα             |
| 10 | 4.4 ± 0.4     | 0 ± 0.3       | 0 ± 0.8                     | Lyα, Hα         |
| 11 | 5.8 ± 0.4     | 2 ± 0.1       | 3 ± 0.1                     | Lyα, Hα         |
| 12 | 2.2 ± 0.6     | 29 ± 4        | 111 ± 22                    | Lyα, Hα, BB     |
| 13 | 0.0 ± 0.4     | 6 ± 0.2       | 22 ± 9                      | BB              |
| 14 | 0.4 ± 0.7     | 16 ± 2        | 5 ± 15                      | BB              |
| 15 | 5.2 ± 0.4     | 2 ± 0.5       | 25 ± 15                     | Lyα             |
| 16 | 0.7 ± 0.5     | 50 ± 20       | 115 ± 32                    | BB              |
| 17 | 1.8 ± 0.5     | 2 ± 0.2       | 6 ± 0.4                     | Morphology      |
| 18 | 1.8 ± 0.4     | 2 ± 0.2       | 9 ± 0.8                     | Morphology      |
| 19 | 1.3 ± 0.5     | 0 ± 0.1       | 0 ± 0.2                     | Morphology      |

1 http://www.mpa-garching.mpg.de/millennium/
Figure 4. Comparison between the derived stellar masses of the observed galaxies within 150 kpc of the Spiderweb galaxy (black histograms) and the masses of galaxies within the same projected distance from the 10 most massive galaxies within 2 SAMs (F08 – blue, and DBL07 – red). Poisson errors for the SAMs are given by $\sim (\text{number in each bin})^{1/2} \times 0.18$. Only those simulated galaxies which have $M_\text{H}$ less that the 3$\sigma$ detection limit (26.6 mag) are plotted. Derived upper mass limits are plotted as black solid circles, and lower mass limits for the satellites are plotted as open circles.

We compile catalogues of the SDSS $g'$, SDSS $i'$, $H$ and $K$ magnitudes for all simulated galaxies. Although the filter curves differ slightly to the observed $g_{275}$, $i_{814}$, $H_{160}$ and $K_s$ filters, this will not affect our comparisons as we only consider large (1 mag) bins when comparing galaxy magnitudes. We convert the SDSS $g'$ luminosities to star formation rates using equation (1) for comparison to the observed UV star formation rates. The SAMs predict a large number of low luminosity satellite galaxies that may bias any comparison to the data, therefore we make a luminosity cut equivalent to the 3$\sigma$ detection limit in the $H_{160}$.

The average number of galaxies within each projected box is 27 for the F08 models and 30 for DBL07 compared to the 19 galaxies observed within 150 kpc of MRC 1138-262. The SAMs predict too many satellite galaxies above our detection threshold compared to our observations. This discrepancy is further compounded by the fact that the dark matter halo of MRC 1138-262 is probably more massive than its simulated counterparts in the SAMs. The missing simulated satellites are faint and have low masses (see Fig. 4). Such faint, low-mass galaxies may not be resolved in the smoothed images, especially if they are not compact. A larger number of galaxies are seen in the higher resolution $HST I_{814}$ images of Miley et al. (2006) of the same region, and the intracluster light discussed in Hatch et al. (2008) may be partly due to unresolved, low mass satellites. It is therefore likely that the observations presented here miss some low-mass satellite galaxies because they cannot be resolved, and thus the discrepancy is most likely due to missing satellite galaxies in the observations, rather than the simulations predicting too many satellite galaxies.

5.2.1 Stellar Masses

The stellar mass of the most massive semi-analytic galaxy at $z = 2.2$ is $2.2 \times 10^{11} M_\odot$ in the F08 models and $2.8 \times 10^{11} M_\odot$ in the DLB07 models. These are $\sim 4$ times less massive than the radio galaxy (galaxy 1). Our derived mass of $10^{12} M_\odot$ for the radio galaxy places it as one of the most massive galaxies detected at $z > 2$. However, we note that this mass is likely to be an upper limit, as we have not taken into account light scattered from the central AGN. Mass derived from fitting stellar templates to broadband photometry are systematically overestimated by a factor of 1.4 – 2 compared to masses derived from fitting stellar templates to spectroscopy (Shapley et al. 2005, Kriek et al. 2008). We must also consider that stellar population synthesis modeling contains large systematic uncertainties and the mass of these galaxies, especially galaxy 1, may be lower than the presented in Table 1. Most high redshift radio galaxies at $z \sim 2$ have similar masses to the most massive galaxies in the SAMs (Seymour et al. 2007), the Spiderweb Galaxy has an exceptionally high mass.

In general the central galaxy in each model contains between 5 and 35 per cent of the stellar mass within 150 kpc of the massive galaxy. If we assume the upper mass limits for the satellite galaxies, the central galaxy in the Spiderweb system contains 60$\pm$13 per cent of the stellar mass. Although the stellar mass distribution of the SAMs is generally not as centrally concentrated as the Spiderweb system, one of the massive galaxies in the F08 SAM contains 35 per cent of the stellar mass within 150 kpc of the massive galaxy (galaxy 1). Our derived mass of $10^{10.5} M_\odot$, which is smaller than those from dark-matter only simulations, and their dominant BCGs contains twice as much stellar mass as any satellite galaxy. These simulations appear to match the derived mass distribution of the Spiderweb system better than the SAMs based on dark-matter only.

Fig 4 compares the derived upper and lower limits of the stellar masses of the galaxies within 150 kpc of the radio galaxy to those galaxies within 150 kpc of the 10 most massive galaxies in the SAMs at $z = 2.2$. Filled circles plot the upper limit masses, whilst open circles plot the lower limit masses. We only plot simulated galaxies with $m_H < 26.6$, the 3$\sigma$ detection limit in $H_{160}$.

The stellar masses of the simulated central galaxies are at least a factor of 4 less that the observed radio galaxy, therefore we are comparing the observations to the masses of galaxies in lower mass halos that of the Spiderweb Galaxy. We thus expect the true distribution of galaxy masses to be shifted to higher masses compared to these simulations. Above $10^{10} M_\odot$, the SAMs agree with the derived lower mass limits, whilst the upper mass limits lie to the right of the simulated galaxy mass distribution. This is entirely consistent with our expectations and imply that the derived upper and lower mass limits are likely to bracket the true stellar mass of the galaxies. Below $10^{10} M_\odot$, the simulations predict a larger number of satellites than are observed. Comparing our upper mass limits to the observed $K_s$ magnitude, we find we are only complete to masses above $\sim 10^{9.8} M_\odot$, thus we are likely to be missing many of the low-mass satellites with red colours.

5.2.2 Galaxy colours

The right-most panel of Fig 5 shows a comparison of the range of UV continuum slope, $\beta$, for the observed and simulated galaxies within 150 kpc of the radio galaxy. The simulated galaxies have been subjected to the same select criteria as the observations – the galaxies must be observed in both the $i'$ and $g'$ bands – thus the galaxies must be actively forming stars. The red UV colours suggest significant amounts of dust between us and the galaxies. The amount of extinction, $E(B-V)$, that can cause these red colours can range from 0.1 to 0.5 depending on the age and metallicity of the

© RAS, MNRAS 000, 000-000
galaxy’s stellar population. DLB07 models provide the closest fit to the data, however, the simulated galaxies with active star formation have too little dust or the stellar ages are too low.

5.2.3 Star formation rates

The SAMs predict total star formation rates of 30–70 M$_\odot$ yr$^{-1}$ within 150 kpc radius, similar to the observed UV star formation rate of 63±8 M$_\odot$ yr$^{-1}$ (dust-uncorrected) in the Spiderweb system. The galaxy colours show that the dust content of the simulated galaxies is less than that of the observed galaxies, and therefore a comparison of the intrinsic star formation rates may result in a discrepancy.

We make a more detailed comparison by comparing the observed $K$ and $g$ magnitudes with those from the SAMs in Fig. 5. The left panel shows there is general agreement between the observed $K$s luminosity function and the SAMs. The $K$s band is the reddest observational band and covers rest-frame $V$–band light. From the selection of observational data used in this work, it is the best tracer of galaxy mass, however there can be significant contamination of light from young stars.

The $g$–band covers rest-frame 1500Å light, which can be converted to star formation rate through equation[1]. This comparison is shown in the middle panel of Fig. 5. We note that the observed $g_{475}$ luminosity function does not have the usual shape of a UV luminosity function. Therefore we do not expect the simulations, which reproduce the general galaxy population at $z \sim 2$ (Guo & White 2008), to match the observed rest-frame UV luminosity function. Indeed the general shape of the simulated UV luminosity function agrees with studies of the $z \sim 2.2$ luminosity function (e.g., Sawicki & Thompson 2006). However this means the simulated satellite galaxies typically have lower $g$ magnitudes, and therefore lower star formation rates, than the observed satellite galaxies in the Spiderweb system. Since the number of simulated galaxies at a given mass approximately matches the observations, this discrepancy is caused by the low star formation rates of the simulated galaxies, rather than a lack of satellite galaxies. The SAMs predict centrally star formation, whilst the observations show that most of the star formation occurs beyond the central radio galaxy. The SAMs reproduce the general galaxy population at $z \sim 2$, but are unable to match some of the details of this special system.

5.2.4 Smooth-particle hydrodynamical simulations

Saro et al. (2009) simulate two protoclusters using smooth particle hydrodynamical simulations and have compared their simulations to observations of the Spiderweb Galaxy, and the galaxies within 75 kpc of this radio galaxy. The BCGs from these smooth-particle hydrodynamical simulations have stellar masses of $5 \times 10^{12} h^{-1} M_\odot$ and $8 \times 10^{11} h^{-1} M_\odot$ at $z = 2$. However, these simulations do not include AGN feedback, which will act to reduce these masses, thus these masses should be considered upper limits. The most massive satellite galaxies within the simulations are a factor of 4-10 less massive than the already dominant BCG. Therefore the stellar mass distributions of the smooth-particle hydrodynamical simulations are different to that of the SAMs, and they provide a better match to the observed stellar mass distribution of the Spiderweb Galaxy.

The smooth particle hydrodynamical simulations have total star formations rates of 600 M$_\odot$ yr$^{-1}$ and 1750 M$_\odot$ yr$^{-1}$ within a smaller region (within a box of 150 h$^{-1}$kpc on a side) than that considered in the data presented here. These high rates are much greater than observed in the UV, although the bright sub-mm emission from this source (Stevens et al. 2003) indicates a large amount of obscured star formation of the order of a thousand solar masses per year. Thus the Saro et al. (2009) results are consistent with the data.

A comparison to the luminosities of the simulated galaxies of

---

Figure 5. $K$ and $G$ band luminosity functions and UV slope for the observed galaxies (black) and 2 versions of SAMs (F08 – blue, and DLB07 – red). Poisson errors for the SAMs are given by $\sim$(number in each bin)$^{1/2} \times 0.18$. Both luminosity functions stop at the 3$\sigma$ detection limit in the considered band. The observed $K$ luminosity function matches reasonably well with the models, suggesting that the mass of the simulated galaxies matches the observations. However there is a mismatch in the observed $g_{475}$ and $G$ luminosity functions implying the star formation rates in the simulated galaxies are too low. Both models produce galaxy UV slopes that are too blue, although the DLB07 models are a better match to the observations. This implies the simulated galaxies which have active star formation have either too little dust or the stellar ages are too low.
Saro et al. (2009) is not possible, but we note that they match the $I_{814}$ galaxy luminosity function of the galaxies within 75 kpc of MRC 1138-262 once they include a prescription for dust attenuation.

6 DISCUSSION

6.1 The state of play at $z = 2.2$

Within a 150 kpc radius, the high redshift radio galaxy MRC1138-262 at $z = 2.2$ is surrounded by 18 satellite galaxies with masses between $10^8 - 10^{11} \, M_{\odot}$. Most of the galaxies within 150 kpc of the radio galaxy are UV-bright line-emitters and have large star formation rates. These star forming galaxies have similar upper limit masses to the galaxies selected through their strong Balmer breaks. Thus the galaxies with sufficient gas reservoirs to fuel large star formation rates do not necessarily have low stellar masses. It is known that star formation can be triggered by merging (e.g. Mihos & Hernquist 1994) and it is likely that this is occurring on a large scale in this system, with most ($> 75 \pm 2$ per cent) of the stars forming in the satellite galaxies, rather than in the central galaxy.

With a mass of $10^{12} \, M_{\odot}$ the central galaxy is already of a similar mass to cD galaxies at the centre of nearby galaxy clusters. The fraction of mass in the Spiderweb system that is contained within the satellite galaxies, depends on whether we assume the stellar masses are better estimated by the lower or upper mass limits, given in columns 3 and 4 of Table 1 respectively. Assuming upper (lower) mass limits, the combined mass within all the satellite galaxies comprises 40 ± 13 (15 ± 5) per cent of the total stellar mass. The radio galaxy still has room for further growth, and may undergo a number of major-mergers before it reaches the limit of the most massive, local cD galaxies. However, galaxy 1 has a $K_s$-band magnitude that is 2 magnitudes brighter than any other protocluster galaxy candidate known within a 1.5 Mpc radius (Kurk et al. 2004), and therefore it is unlikely that the protocluster contains more galaxies of a similar mass at the observed redshift.

De Lucia & Blaizot (2007) find that ~70 per cent of BCGs at $z = 2$ do not become BCGs at $z = 0$, but rather end up as the central galaxies of massive halos which are less massive than cluster halos. The most massive halos at $z = 2$ are generally not the progenitors of $z = 0$ clusters. Therefore, whilst MRC 1138-262 appears to be the brightest protocluster galaxy at $z = 2$, it is uncertain whether the Spiderweb system will evolve into a local BCG.

6.2 The evolution of the Spiderweb System

6.2.1 Merging timescales

A complete N-body numerical investigation of this system is required to study the evolution of the Spiderweb system, which comprises of the radio galaxy and all galaxies within a 150 kpc projected radius. The data given in Table 1 and Fig 2 can be taken as initial conditions for such a simulation. However, if we assume that the satellite galaxies all lie on circular orbits around the radio galaxy, with radii given by their projected radii, we can use analytic approximations to gain an idea of the evolution of this system.

When an orbital point mass (such as the satellite galaxy) travels through an uniform background mass distribution (such as the dark matter halo of the central galaxy), it experiences a steady deceleration known as dynamical friction (Chandrasekhar 1943). The merging timescale ($\tau_{\text{merge}}$) of these galaxies can be taken as the dynamical friction timescale given by

$$\tau_{\text{merge}} = 1.17 \frac{\tau_{\text{dyn}}}{\ln \Lambda} \left( \frac{M_{\text{central}}}{M_{\text{satellite}}} \right),$$

where the Coulomb logarithm $\ln \Lambda = \ln(1 + M_{\text{central}}/M_{\text{satellite}})$, $M_{\text{central}}$ is the mass of the central galaxy, $M_{\text{satellite}}$ is the mass of the satellite galaxies and the orbital timescale is given by $\tau_{\text{dyn}} = (\pi R^{3/2})/\Lambda$, where $R$ is the distance between the satellite and central galaxy (Binney & Tremaine 1987). $M_{\text{central}}/M_{\text{satellite}}$ is the ratio of dark matter halo mass. Matching the abundances of galaxies above a given stellar mass to halos above a given dark-matter mass is a powerful technique employed to relate the galactic stellar mass to a halo mass (Conroy & Wechsler 2008), however, this technique is not reliable at high redshift ($z > 1$). Instead we assume that the amount of dark matter within each galaxy is directly proportional to the galaxy’s stellar mass, therefore the ratio $M_{\text{central}}/M_{\text{satellite}}$ is the ratio of stellar mass. If the abundance-matching estimates of (Conroy & Wechsler 2008) are correct, the merging timescales will be reduced, especially for low mass satellites. $M_{\text{satellite}}$ is the mass of the satellite galaxies and $M_{\text{central}} = 1 \times 10^{12} \, M_{\odot}$, the lower limit of the stellar mass of galaxy 1. The total (dark matter plus stellar) mass of the central galaxy is estimated from the available dynamical information. Spectroscopic observations of galaxies 5, 7, 9, 11, 15 show they have line-of-sight velocities of 1030 km s$^{-1}$, 210 km s$^{-1}$, 570 km s$^{-1}$, 1180 km s$^{-1}$ and 1240 km s$^{-1}$ relative to the central galaxy (Miley et al. 2006) which are consistent with a central mass of $M_{\text{total}} = 10^{13} \, M_{\odot}$.

Boylan-Kolchin et al. (2008) used numerical simulations to find a more realistic merging timescale for minor mergers. The merging timescale is increased relative to equation (2) due to the significant mass loss of the satellites as they orbit the central galaxy. The merging timescale is given by

$$\tau_{\text{merge}} = 0.216 \frac{\tau_{\text{dyn}}}{\ln \Lambda} \left( \frac{M_{\text{central}}}{M_{\text{satellite}}} \right)^{1.3} \exp[1.9(1 - e^2)^{1/2}],$$

where $e$ is the orbital eccentricity. We assume all satellites have circular orbits ($e = 0$). Lower angular momentum orbits have shorter merging times, if the satellites have elliptical orbits ($0 < e < 1$), their merging times could decrease by a factor of ~7.

The merging times for all satellite galaxies within 150 kpc of the central radio galaxy are shown in Fig 6. The filled circles mark the merging timescales derived using equation (3) and the upper limit of the galaxy mass, whilst the open circles use the lower limit of the galaxy mass. The central galaxy is always assumed to have a mass of $1.1 \times 10^{12} \, M_{\odot}$, which is consistent with other measurements from the rest-frame near-infrared. The open circles indicate the longest possible merging times, as we have assumed both the lowest masses and orbital eccentricities of the satellite galaxies. The vertical lines show the time from $z = 2.2$ until $z = 1$ and $z = 0$.

The more massive galaxies generally merge earliest, irrespective of their initial radial separation from the central radio galaxy. Most of the satellite galaxies will merge by $z = 0$ if they have masses close to their upper mass limits. In doing so the radio galaxy will increase its mass by a factor of $1.5^{+0.4}_{-0.3}$. Alternatively, the mass of the satellite galaxies may be closer to the lower mass limits. In this case, only a third of the satellites will merge before $z = 0$. Thus the amount of growth is dramatically reduced, with only a ~13 per cent increase in the stellar mass through merging. However, this is not the full story, as tidal stripping is an important process that also depends on the mass of the satellite galaxies. The minimum merging timescale is ~300 Myr which allows the satellites to evolve.
Figure 6. The merging times for all satellite galaxies within 150 kpc of the central radio galaxy. The dotted and dashed lines indicate the time from $z = 2.2$ until $z = 0$ and $z = 0$ respectively. Merging timescales are derived from equation (3) and are from Table 1. Filled circles indicate merging times using the upper mass limits from column 4 of Table 1, whilst the open circles are merging timescales derived assuming the satellites have the lower masses from column 3. Red symbols are galaxies which are likely to be experiencing rapid tidal stripping as their tidal radii are less than their half-light diameters. Most satellites will either merge with the central radio galaxy by $z = 0$, or are experiencing rapid tidal stripping, and are therefore unlikely to survive until $z = 0$ as separate entities.

Figure 7. The tidal radius of the satellite galaxies derived from the upper mass limit of the satellites shown as filled circles, whilst radii derived from lower mass limits are open circles. Red points are galaxies whose tidal radii are smaller than their half-light diameters as measured from the $H_{160}$ image. All of these satellites are likely to be experiencing some tidal stripping, but those coloured red will be stripped at a faster rate.

6.2.2 Tidal stripping

The typical orbital time is 100 Myrs, therefore the satellite galaxies will orbit the central galaxy between several and hundreds of times before they merge. As the satellites orbit the central galaxy, tidal interactions can strip large quantities of gas and stars from the satellite galaxies into tidal tails, which are only partly re-acquired by the galaxies (di Matteo et al. 2007). After several orbits kinematically distinct streams will appear. One possible such stream lies between galaxies 17 and 18 (see Figs. 1 and 2).

The tidal radius of the satellites can be approximated by

$$ R_{\text{tidal}} \approx \left( \frac{M_{\text{satellite}}}{3M_{\text{central}}} \right)^{\frac{1}{2}} R, $$

where $R$ is the distance between the satellite galaxy and the central radio galaxy (Binney & Tremaine 1987). This radius approximately marks the limit within which stars are bound to the satellite, and beyond which stars can be tidally stripped. Fig. 7 shows that the tidal radius for many of the satellite galaxies is less than 10 kpc, hence the outer stars of the satellites are likely being stripped away. Further stripping may occur from interactions between the satellite galaxies, including those galaxies below the observational detection limit. The mass ratios of the satellite galaxies are closer to unity, which increases the tidal radii. However, the radial separation between two satellites may become small during some part of their orbits, and during this period the forces between the satellite galaxies may dominate, increasing the amount of tidal stripping that occurs. Such interactions can only be modeled by a N-body numerical simulation of the system.

The red symbols in Figs. 6 and 7 mark the galaxies which have smaller tidal radii than half-light diameters, measured from the $H_{160}$ image. These galaxies are probably experiencing tidal stripping of their central stars, and it is unlikely that these galaxies will survive as distinct entities for a large number of orbits. It is likely that most of their stars will be stripped before they merge. There is a dependance on mass, so the satellite galaxies would experience more stripping if their mass was closer to the lower mass limit rather than the upper mass limit.

Fig. 7 shows that the stars and gas can be stripped from the satellite galaxies at large distances, more than 50 kpc from the central galaxy. Therefore these stripped stars may form an extended stellar halo around the radio galaxy, so the radio galaxy can grow in size, or form something akin to a cD halo. It has been postulated that the extended stellar halos of cD galaxies may form as stars are stripped during mergers (Gallagher & Ostriker 1972; Malumuth & Richstone 1984) and these observations may have captured the Spiderweb Galaxy in the early stages of such a process. A numerical N-body simulation of this system will determine whether such conditions can result in an extended stellar halo such as a cD halo.

6.2.3 Assembly scenarios of the Spiderweb system

The evolution of the system strongly depends on the mass ratio of the central and satellite galaxies. Since we have assumed that the central galaxy has a mass of $10^{12} M_{\odot}$, consistent with previous measurements, the evolution of this system depends on the mass of the satellite galaxies. If the galaxies have true masses similar to the derived upper mass limits, most of the satellites will merge by $z = 0$, with the most massive systems merging within a few hundred Myrs. Tidal stripping is most important for the low mass satellites that are within 30 kpc of the central galaxy. The stars accumulated from the mergers and stripping are likely to be deposited close to the central
Alternatively, if the masses of the satellite galaxies are similar to the lower limits we have presented, only the third most massive satellites will merge with the central galaxy by $z = 0$, and most merging occurs at later times, between $1 < z < 0$. The mass of the central galaxy is not greatly increased by these mergers, although this is due to the low masses of the satellites, rather than the low number of mergers. The other satellites are unlikely to survive as distinct entities until $z = 0$. Fig[6] shows that these satellites that will not merge by $z = 0$ are experiencing strong tidal stripping of their inner stars. These stars will be stripped from the satellites and become bound to the central galaxy. In such a manner the central galaxy may accumulate all the stellar mass that is currently observed in the satellites. The satellites experiencing strong tidal stripping generally have the lowest mass, and longest merging timescale, and can be stripped at large separations from the central galaxy, up to 60 kpc.

In both scenarios the central galaxy is likely to accumulate most of the stellar mass of the satellite galaxies by $z = 0$, however, it does so in different manners. If the satellites’ masses are large, most of the stellar mass from these satellites is deposited relatively quickly and at small radial distances from the core of the central radio galaxy. On the other hand, if the satellites have low masses, most mass is deposited through relatively slow mergers. A fraction of the stellar mass will be deposited at large radii from the core of the central galaxy by tidal stripping, forming an extended stellar halo. The total growth of the central galaxy due to the consumption of the satellite galaxies is between ~10 and 120 per cent. This large range is due to the large uncertainty of the stellar mass derived through stellar template synthesis modeling. The future evolution of this system is likely to be somewhere in between the two scenarios present above. In addition, this amount of growth should be considered a lower limit, as massive nearby galaxies ($> 10^{11} M_\odot$), that lie beyond 150 kpc, may also merge with the central galaxy.

The crucial parameter that drives the evolution of such a system is the mass ratio of the central and satellite galaxies. If a galaxy is able to assemble quickly such that it is times many more massive than the general galaxy population at high redshift ($z > 2$), a larger fraction of the stellar mass from its future mergers and close encounters will end up in an extended stellar halo. Thus whether or not a brightest cluster galaxy becomes a cD galaxy, may depend on its early-time formation history. The Spiderweb galaxy, with its large mass at $z = 2.2$, is in a good position to build a substantial extended stellar halo by $z = 0$.

7 CONCLUDING REMARKS

We have used multi-wavelength high-resolution images to study the assembling radio galaxy, MRC1138-262, at $z \sim 2.2$. We identified a galaxy population that lies at the same redshift and within a projected distance of 150 kpc from the radio galaxy. Fitting stellar population models to the photometry of these galaxies we have examined the current distribution of stellar mass and ongoing star formation, and used analytic approximations to predict the future evolution of this galaxy.

1. We identified 19 protocluster galaxies, comprising of a central massive galaxy ($10^{12} M_\odot$) that is surrounded by 18 smaller satellite galaxies with masses between $10^8$ and $10^{11} M_\odot$. The central galaxy is already of a similar mass to local cD galaxies, but there is still room for growth as cD galaxies can have masses of $10^{13} M_\odot$. The satellite galaxies contribute between 10 and 50 per cent of the stellar mass of the whole system, and they dominate the star formation rate, containing 75 ± 2 per cent of the instantaneous star formation (dust-uncorrected).

2. The majority of the satellite galaxies have large star formation rates and are inferred to be gas-rich. The intrinsic star formation rates are uncertain because the red UV continuum slopes indicate dust obscuration. The high star formation rates of the satellite galaxies coupled with the presence of dust explains the spatially extended 850µm emission from MRC 1138-262 [Seyfert et al 2003]. The extended sub-mm emission most likely results from the sum of the compact star formation occurring in the radio and satellite galaxies rather than a system-wide starburst. These results predict that a large fraction of the sub-mm flux should come from the west of the radio galaxy and high resolution sub-mm observations may be able to separate the emission into the various galaxy components.

3. We have compared the properties of the galaxies within 150 kpc of MRC 1138-262 to predictions of semi-analytic models. The models are able to reproduce the observed K-band galaxy luminosity function, and the derived stellar masses of the satellite galaxies. The models predict galaxy colours which are bluer than the observations, and yet they do not match the observed high star formation rates. There may be a process, not accounted for in the models, which increases star formation when galaxies are situated in a dense merging environment, such as the Spiderweb system.

4. Most of the satellite galaxies within 150 kpc will accrete onto the central radio galaxy before $z = 0$, increasing its mass by between ~10 and 120 per cent. The merging timescales are long, and the vast majority of the galaxies will not merge for at least 300 Myrs, by which time they will have completed many orbits. Gas is rapidly converted to stars as inferred from the large star formation rates. Therefore the final merging galaxies are likely to contain very little gas, and the mergers are likely to be dissipationless.

5. The tidal radii of the satellite galaxies are small, generally less than 10 kpc, and in some cases are less than the half-light diameter of the galaxies. Therefore stars and gas are probably being stripped from the satellite galaxies. This stripping can occur at large distances from the central galaxy (more than 60 kpc) and the stripped stars may form an extended stellar halo such as those that encompass local cD galaxies.

8 ACKNOWLEDGMENTS

We would like to thank the anonymous referee for helping to improve and clarify the paper, Ian McCarthy for help with the Millennium simulations, Rychard Bouwens for help with the reduction of the NICMOS data, and Alex Saro for useful discussions. This research has been based on observations made with the NASA/ESA Hubble Space Telescope, obtained at the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5-26555. These observations are associated with programs 10327 and 10404. The Millennium Simulation databases used in this paper and the web application providing online access to them were constructed as part of the activities of the German Astrophysical Virtual Observatory. NAH and GM acknowledge funding from the Royal Netherlands Academy of Arts and Sciences. JK is supported by the DFG, Sonderforschungsbereich (SFB) 439.
REFERENCES

Benitez N., 2000, ApJ, 536, 571
Bertin E., Arnouts S., 1996, A&A, 117, 393
Binney J., Tremaine S., 1987, Galactic dynamics. Princeton, NJ, Princeton University Press, 1987, 747 p.
Blakeslee J. P., Anderson K. R., Meurer G. R., Benitez N., Magee D., 2003, in Payne H. E., Jedrzejewski R. I., Hook R. N., eds, Astronomical Data Analysis Software and Systems XII Vol. 295 of Astronomical Society of the Pacific Conference Series, An Automatic Image Reduction Pipeline for the Advanced Camera for Surveys. pp 257–+
Boylan-Kolchin M., Ma C.-P., Quataert E., 2008, MNRAS, 383, 93
Bruzual G., Charlot S., 2003, MNRAS, 344, 1000
Chabrier G., 2003, PASP, 115, 763
Chandraseker S., 1943, ApJ, 97, 255
Coe D., Benitez N., Sanchez S. F., Jee M., Bouwens R., Ford H., 2006, AJ, 132, 926
Conroy C., Gunn J. E., White M., 2008, ArXiv e-prints
Conroy C., Wechsler R. H., 2008, ArXiv e-prints
De Lucia G., Blaizot J., 2007, MNRAS, 375, 2
Font A. S., Bower R. G., McCarthy I. G., Benson A. J., Frenk C. S., Helly J. C., Lacey C. G., Baugh C. M., Cole S., 2008, MNRAS, 389, 1619
Gallagher III J. S., Ostriker J. P., 1972, AJ, 77, 288
Guo Q., White S. D. M., 2008, ArXiv e-prints
Hatch N. A., Overzier R. A., Röttgering H. J. A., Kurk J. D., Miley G. K., 2008, MNRAS, 383, 931
Ivison R. J., Morrison G. E., Biggs A. D., Smail I., Willner S. P., Gurwell M. A., Greve T. R., Stevens J. A., Ashby M. L. N., 2008, ArXiv e-prints, 808
Kodama T., Tanaka I., Kajisawa M., Kurk J., Venemans B., De Breuck C., Vernet J., Lidman C., 2007, MNRAS, 377, 1717
Kriek M., van Dokkum P. G., Franx M., Illingworth G. D., Marchesini D., Quadri R., Rudnick G., Taylor E. N., Förster Schreiber N. M., Gawiser E., Labbé I., Lira P.,uyts S., 2008, ApJ, 677, 219
Kurk J. D., 2003, PhD thesis, Leiden University, P.O. Box 9504, 2300 RA Leiden, The Netherlands
Kurk J. D., Pentericci L., Röttgering H. J. A., Miley G. K., 2004, A&A, 428, 793
Malumuth E. M., Richstone D. O., 1984, ApJ, 276, 413
Maraston C., Daddi E., Renzini A., Cimatti A., Dickinson M., Papovich C., Pasquali A., Pirzkal N., 2006, ApJ, 652, 85
Mihos J. C., Hernquist L., 1994, ApJ, 425, L13
Miley G., De Breuck C., 2008, ARA&A, pp 1–+
Miley G. K., Overzier R. A., Zirm A. W., Ford H. C., Kurk J., Pentericci L., Blakeslee J. P., Franx M., Illingworth G. D., Postman M., Rosati P., Röttgering H. J. A., Venemans B. P., Helder E., 2006, ApJ, 650, L29
Overzier R. A., Miley G. K., Ford H. C., 2007, New Astronomy Review, 51, 202
Papovich C., Dickinson M., Ferguson H. C., 2001, ApJ, 559, 620
Pentericci L., Kurk J. D., Röttgering H. J. A., Miley G. K., van Breugel W., Carilli C. L., Ford H., Heckman T., McCarthy P., Moorwood A., 2000, A&A, 361, L25
Pentericci L., Roettgering H. J. A., Miley G. K., Carilli C. L., McCarthy P., 1997, A&A, 326, 580
Reddy N. A., Erb D. K., Steidel C. C., Shapley A. E., Adelberger K. L., Pettini M., 2005, ApJ, 633, 748
Salim S., Rich R. M., Charlot S., Brinchmann J., Johnson B. D., Schiminovich D., Seibert M., Mallery R., Heckman T. M., Forster K., Friedman P. G., 2007, ApJS, 173, 267
Saro A., Borgan S., Tornatore L., De Lucia G., Dolag K., Murante G., 2009, MNRAS, 392, 795
Saro A., De Lucia G., Dolag K., Borgan S., 2008, ArXiv e-prints
Sawicki M., Thompson D., 2006, ApJ, 642, 653
Schlegel D. J., Finkbeiner D. P., Davis M., 1998, ApJ, 500, 525
Seymour N., Stern D., De Breuck C., Vernet J., Rettura A., Dickinson M., Dey A., Eisenhardt P., Fosbury R., Lacy M., McCarthy P., Miley G., Rocca-Volmerange B., Rottgering H., Stanford S. A., Teplitz H., van Breugel W., 2007, ArXiv Astrophysics e-prints
Shapley A. E., Steidel C. C., Erb D. K., Reddy N. A., Adelberger K. L., Pettini M., Barnby P., Huang J., 2005, ApJ, 626, 698
Spergel D. N., Verde L., Peiris H. V., Komatsu E., Nolta M. R., Bennett C. L., Halpern M., Hinshaw G., Jarosik N., Kogut A., Limon M., Meyer S. S., Page L., Tucker G. S., Weiland J. L., Wollack E., Wright E. L., 2003, ApJS, 148, 175
Stevens J. A., Ivison R. J., Dunlop J. S., Smail I. R., Percival W. J., Hughes D. H., Röttgering H. J. A., van Breugel 2003, Nature, 425, 264
Venemans B. P., Röttgering H. J. A., Miley G. K., van Breugel W. J. M., de Breuck C., Kurk J. D., Pentericci L., Stanford S. A., Overzier R. A., Croft S., Ford H., 2007, A&A, 461, 823
Villar-Martín M., Humphrey A., De Breuck C., Fosbury R., Binette L., Vernet J., 2007, MNRAS, 375, 1299
Villar-Martín M., Tadhunter C., Clark N., 1997, A&A, 323, 21
Zirm A. W., Overzier R. A., Miley G. K., Blakeslee J. P., Clampin M., De Breuck C., Demarco R., Ford H. C., et al. 2005, ApJ, 630, 68
Zirm A. W., Stanford S. A., Postman M., Overzier R. A., Blakeslee J. P., Rosati P., Kurk J., Pentericci L., Venemans B., Miley G. K., Röttgering H. J. A., Franx M., van der Wel A., Demarco R., van Breugel W., 2008, ApJ, 680, 224