HAT-P-5b: A JUPITER-LIKE HOT JUPITER TRANSITING A BRIGHT STAR

G. Á. Bakos, A. Shporer, A. Pál, G. Torres, Géza Kovács, D. W. Latham, T. Mazeh, A. Ofir, R. W. Noyes, D. D. Sasselov, F. Bouchy, F. Pont, D. Queloz, S. Udry, G. Esquerdo, B. Sipőcz, Gábor Kovács, R. Stefanik, J. Lázár, I. Papp, and P. Sári

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ABSTRACT

We report the discovery of a planet transiting a moderately bright ($V = 12.00$) G star, with an orbital period of 2.788491 ± 0.000025 days. From the transit light curve we determine that the radius of the planet is $R_p = 1.257 ± 0.053 R_J$, HAT-P-5b has a mass of $M_p = 1.06 ± 0.11 M_J$, similar to the average mass of previously known transiting exoplanets, and a density of $ρ_p = 0.66 ± 0.11$ g cm$^{-3}$. We find that the center of transit is $T_c = 2,454,241.77663 ± 0.00022$ days (HJD), and the total transit duration is 0.1217 ± 0.0012 days.

Subject headings: planetary systems — stars: individual (GSC 02634–01087, HAT-P-5)

1. INTRODUCTION

To date about 20 extrasolar planets have been found that transit their parent stars and thus yield values for their mass and radius. Masses range from 0.07 $M_J$ (GI 436; Gillon et al. 2007) to about 9 $M_J$ (HAT-P-2b; Bakos et al. 2007), and radii from 0.4 $R_J$ (GI 436) to about 1.7 $R_J$ (TRES-4; Mandushev et al. 2007). These data provide an opportunity to compare observations with theoretical models of planetary structure across a wide range of parameters, including those of the host star (e.g., Burrows et al. 2007; Fortney et al. 2007 and references therein). Transits also yield precise determination of other physical parameters of the extrasolar planets, for instance the surface gravity. Interesting correlations between these parameters were noted early on, such as that between masses and periods (Mazeh et al. 2005) or periods and surface gravities (Southworth et al. 2007). Classes of these close-in planets have also been suggested, such as very hot Jupiters (VHJs; $P = 1–3$ days) and hot Jupiters (HJs; $P = 3–9$ days; Gaudi et al. 2005), or a possible dichotomy based on Safronov numbers (Hansen & Barman 2007). However, the small ensemble of transiting exoplanets (TEPs) does not allow robust conclusions; thus the addition of new discoveries is valuable.

Over the past year the HATNet project (Bakos et al. 2002, 2004), a wide-angle photometric survey, has announced four TEPs. In this Letter we report on the detection of a new transiting exoplanet, which we label HAT-P-5b, and our determination of its parameters, such as mass, radius, density, and surface gravity.

2. OBSERVATIONS AND ANALYSIS

2.1. Detection of the Transit in the HATNet Data

GSC 02634–01087, also known as 2MASS J18173731+3637170, is a G star with $I = 11.3$ and $V = 12.00$. It was initially identified as a transit candidate in our internally labeled field G196, centered at $α = 18°08′$, $δ = 37′′0′′$. The data were acquired by HATNet’s 7-m telescope at the Fred Lawrence Whipple Observatory (FLWO) of the Smithsonian Astrophysical Observatory (SAO) and HAT-9 telescope at the Submillimeter Array (SMA) site atop Mauna Kea, Hawaii. Following a standard calibration procedure of the frames (meaning bias, dark, and flat-field corrections), data were reduced using the astrometry code of Pál & Bakos (2006) and a highly fine-tuned aperture photometry. We applied our external parameter decorrelation technique on the light curves, whereby deviations from the median were cross-correlated with a number of “external parameters,” such as the $X$ and $Y$ subpixel position, FWHM, hour angle, and zenith distance. We have also applied the trend filtering algorithm (Kovács et al. 2005), along with the box least squares (BLS; Kovács et al. 2002) transit-search algorithm in our analysis. For field G196 we gathered ∼3750 (HAT-7) plus ∼890 (HAT-9) data points at 5.5 minute cadence between 2005 June 8 and 2005 December 5 (UT). In the light curve of star GSC 02634–01087 we detected a ∼13 mmag transit with a 2.7881 day period, a signal-to-noise ratio of 12 in the BLS frequency spectrum, and a dip significance of 18 (Kovács & Bakos 2005). The top panel of Figure 1 shows the unbinned light curve with all ∼4640 data points, folded with the period that we derived subsequently, based on high-precision follow-up photometry, as described in § 2.4 below.

2.2. Early Spectroscopy Follow-Up

Initial follow-up observations were made with the CfA Digital Speedometer (DS; Latham 1992) in order to characterize the host star and to reject obvious astrophysical false-positive scenarios that mimic planetary transits. The four radial velocity (RV) measurements obtained over an interval of 33 days showed an rms residual of 0.41 km s$^{-1}$, consistent with no detectable RV variation. Atmospheric parameters for the star (effective temperature $T_{\text{eff}}$, surface gravity log $g$, metallicity [Fe/H], and projected rotational velocity $v$ sin $i$) were derived as described by Torres et al. (2002). The first three quantities are strongly correlated and difficult to determine simulta-
neously. For example, the unconstrained value \( \log g = 4.0 \pm 0.2 \) we obtained is somewhat lower than derived from our stellar evolution modeling in § 3, which is \( \log g = 4.37 \). Consequently, in a second iteration we held \( \log g \) fixed at this value and redetermined the other quantities, obtaining \( T_{\text{eff}} = 5960 \pm 100 \) K, \( [\text{Fe/H}] = +0.24 \pm 0.15 \), and \( v \sin i = 2.6 \pm 1.5 \) km s\(^{-1}\). These correspond to a slowly rotating early G main-sequence star.

2.3. High-Precision Spectroscopy Follow-Up

High-resolution spectroscopic follow-up was carried out at the Haute Provence Observatory (OHP) 1.93 m telescope, with the SOPHIE\(^{12}\) spectrograph (Bouchy et al. 2006). SOPHIE is a multiorder echelle spectrograph fed through two fibers, one of which is used for starlight and the other for sky background or a wavelength calibration lamp. The instrument is entirely computer-controlled, and a standard data reduction pipeline automatically processes the data on CCD readout. RVs are calculated by numerical cross-correlation with a high-resolution observed spectral template of a G2 star. Similar spectroscopic follow-up with SOPHIE has already resulted in the confirmation of two TEPs: WASP-1b and WASP-2b (Cameron et al. 2007). HAT-P-5 was observed with SOPHIE in the high-efficiency mode (\( R \sim 39,000 \)) during our 2007 May 2–13 observing run, with an additional measurement taken on June 4. Depending on observing conditions, exposure times were in the range of 15–35 minutes, resulting in signal-to-noise ratios of 20–55 pixel\(^{-1}\) at \( \lambda = 5500 \) Å. Using the empirical relation of Cameron et al. (2007) we estimated the RV photon-noise uncertainties to be 10–25 m s\(^{-1}\).

In order to better characterize the transit parameters and also to improve the ephemerides, we performed follow-up photometric observations with 1 m class telescopes. A partial transit of HAT-P-5b was observed using the KeplerCam detector on the FLWO 1.2 m telescope (see Holman et al. 2007) on UT 2007 May 18. We refer to this event as having transit number \( N_\text{tr} = -1 \). Three days later a full transit, \( N_\text{tr} = 0 \), was observed with the same instrument. The two Sloan \( z \)-band light curves are shown in the middle panel of Figure 1. We also gathered data for four subsequent full transit events, \( N_\text{tr} = 2, N_\text{tr} = 11, N_\text{tr} = 16, \) and \( N_\text{tr} = 21 \), using the Wise 1 m telescope in the Cousins R band (bottom panel of Fig. 1). Data were reduced in a manner similar to the HATNet data, using aperture photometry and an ensemble of \( \sim 300 \) comparison stars in the field. Since the follow-up observations span 22 transit cycles (\( \sim 2 \) month time span), we were able to obtain an accurate ephemeris. An analytic model was fit to these data, as described below in § 6, and yielded a period of \( \Delta T = 0.0012 \pm 0.000025 \) days and reference epoch of midtransit \( T_0 = 2.454,241,776.63 \pm 0.00022 \) days (HJD). The length of the transit as determined from this joint fit is \( 0.1217 \pm 0.0012 \) days (2 hr, 55 minutes), the length of ingress is \( 0.0145 \pm 0.0007 \) days (20.9 minutes), and the central transit depth is 0.0136 mag.

3. STELLAR PARAMETERS

The mass (\( M_\star \)) and radius (\( R_\star \)) of a transiting planet scale with those of the parent star. In order to determine the stellar properties needed to place \( M_\star \) and \( R_\star \) on an absolute scale, we made use of stellar evolution models along with the observational constraints from spectroscopy. Because of its relative faintness, the host star does not have a parallax measurement from Hipparcos, and thus a direct estimate of the absolute magnitude is not available for use as a constraint. An alternative approach is to use the surface gravity of the star, which is a sensitive measure of the evolutionary state of the star and therefore has a very strong influence on the radius. However, \( \log g \) is a notoriously difficult quantity to measure spectroscopically and is often strongly correlated with other spectroscopic parameters (see § 2.2). It has been pointed out by Sozzetti et al. (2007) that the normalized separation of the planet,
The noise of 14.4 m s\(^{-1}\), which we attribute to "stellar jitter." This level of jitter is consistent with the predictions of Saar et al. (1998) for a projected rotational velocity such as what we expect for the parent star. The final fit, with the internal errors added in quadrature of uncorrelated velocities (see, e.g., Queloz et al. 2001; Torres et al. 2005). As discussed later in § 6, an analytic fit to the light curve yields \(a/R_\star \approx 7.50 \pm 0.19\).

This value, along with \(T_{\text{eff}}\) and [Fe/H] from § 6, was compared with the Yonsei-Yale stellar evolution models of Yi et al. (2001) following Sozzetti et al. (2007). As described earlier, the initial temperature and metallicity from our DS spectroscopy were subsequently improved by applying the log \(g\) constraint from the models, and repeating the isochrone comparison. This resulted in final values for the stellar mass and radius of \(M_\star = 1.160 \pm 0.062 \, M_\odot\) and \(R_\star = 1.167 \pm 0.049 \, R_\odot\), and an estimated age of 2.6 \pm 1.8 Gyr. We summarize these and other properties in Table 2.

### 4. Spectroscopic Orbital Solution

Our eight RV measurements from SOPHIE were fitted with a Keplerian orbit model solving for the velocity semi-amplitude \(K\) and the center-of-mass velocity \(\gamma\), holding the period and transit epoch fixed at the well-determined values from photometry. The eccentricity was initially set to zero. The resulting rms residual of \(\sim 23.7\, \text{m s}\(^{-1}\) is somewhat larger than expected from the internal errors, and we find that a reduced \(\chi^2\) value of unity necessitates the addition in quadrature of uncorrelated noise of 14.4 m s\(^{-1}\), which we attribute to "stellar jitter." This level of jitter is consistent with the predictions of Saar et al. (1998) for a projected rotational velocity such as what we measure for the parent star. The final fit, with the internal errors increased as described above, yields \(K = 138 \pm 14\, \text{m s}\(^{-1}\) and \(\gamma = 7613.8 \pm 9.1\, \text{m s}\(^{-1}\). The observations and fitted RV curve are displayed in the top panel of Figure 2, with the residuals shown in the middle panel.

As a test we allowed for the possibility of an eccentric orbit and solved for the two additional quantities \(e \cos \omega\) and \(e \sin \omega\), but the results were insignificantly different from zero.

### 5. EXCLUDING BLEND SCENARIOS

We have tested the reality of the velocity variations by carefully examining the spectral line bisectors of our OHP data. If the velocity changes measured are due only to distortions in the line profiles arising from contamination of the spectrum by the presence of a binary with a period of 2.79 days, we would expect the bisector spans (which measure line asymmetry) to vary with this period and with an amplitude similar to the velocities (see, e.g., Queloz et al. 2001; Torres et al. 2005). As shown in the bottom panel of Figure 2, the changes in the bisector spans are of the same order as the residual RV variations, and much smaller than the radial velocity semi-amplitude itself. This analysis shows that the orbiting body is a planet and rules out a possible blend scenario.

### 6. PLANETARY PARAMETERS

The light-curve parameters of HAT-P-5b were determined from a joint fit based on the six distinct transit events, observed with the FLWO 1.2 m and Wise 1 m telescopes. A circular orbit was assumed, based on our analysis above. We adopted a quadratic limb-darkening law for the star, and took the appropriate coefficients from Claret (2004) for both the Sloan \(z\) and Cousins \(R\) bands. The drop in flux in the light curves was modeled with the formalism of Mandel & Agol (2002) using the equations for the general case (i.e., \(not\) the small-planet approximation). The adjusted parameters in the fit were (1) the midtransit times of the first full transit \((N_{\text{tr}} = 0, T_{\text{c1}})\) and the last full transit \((N_{\text{tr}} = 21, T_{\text{c2}})\) (this is equivalent to fitting for an epoch \(E\) and a period \(P\)); (2) the relative planetary radius, \(R_p/R_\star\); (3) the square of the normalized impact parameter, \(b^2\); and (4) the quantity \(\gamma R_\star \equiv a/R_\star (2\pi P)/(1 - b^2)^{1/2}\). From simple geometric considerations, \(\gamma R_\star\) and \(b^2\) have an uncorrelated posteriori probability distribution in parameter space. This amounts to an orthogonalization of the fitted parameters, similar but simpler than the one employed by Burke et al. (2007) for the case of XO-2b.

We used the Markov Chain Monte Carlo algorithm (see, e.g., Holman et al. 2007) to derive the best-fit parameters. Uncertainties were estimated using synthetic data sets, by adding Gaussian noise to the fitted curve at the dates of our observations and resolving the light-curve fit. The magnitude of the noise was taken from the white- and red-noise estimations based on the measured residuals. This process was repeated \(1.5 \times 10^5\) times, yielding a good representation of the a posteriori distribution of the best-fit parameter values. We found this method of error estimation to be robust, since it is not sensitive to the number of out-of-transit points.
The result for the radius ratio is $R_P/R_*=0.1106 \pm 0.0006$, and the normalized separation is $a/R_*=7.50 \pm 0.19$. We found that the a posteriori distribution of $b^2$ is consistent with a symmetric Gaussian distribution and yields $b^2=0.181 \pm 0.040$; therefore the orbit is inclined. From the inclination, the mass of the star (Table 2), and the orbital parameters (§ 4), the other planetary parameters (such as mass and radius) are derived in a straightforward way, and are summarized in Table 3. We note that $a/R_*$, as derived from the light-curve fit, is an important constraint in the stellar parameter determination (§ 3), which in turn defines the limb-darkening coefficients that are used in the light-curve fit. Thus, after the initial fit to the light curve and the stellar parameter determination, we performed another iteration in the light-curve fit. We found that the change in parameters was imperceptible.

The possibility of transit time variations (TTVs) was checked by fitting the center of the transit of the five full transit events independently. We found no sign of TTV, as the transit times differ by less than 1 s from the expected values (listed in Table 3).

7. Conclusions

HAT-P-5b is an ordinary hot Jupiter ($P = 2.788$ days) with slightly inflated radius ($R_p = 1.26 \, R_J$) for its mass of 1.06 $M_J$, orbiting a slightly metal-rich solar-like star. The $\sim 20\%$ radius inflation is what current models predict for a planet with equilibrium temperature of $\sim 1500$ K (Burrows et al. 2007; Fortney et al. 2007).

HAT-P-5b is more massive than any of the known TEPs with similar period (2.5 days $\leq P \leq 3$ days), such as XO-2b, WASP-1, HAT-P-3b, TRES-1, and HAT-P-4b, with the exception of TRES-2. The latter is fairly similar in mass, radius, orbital period, and stellar effective temperature.

However, HAT-P-5b is interesting in that it falls between Class I and II, as defined by the Safronov number and of the planet (Hansen & Barman 2007). HAT-P-5b has a Safronov number of 0.059 $\pm$ 0.005, while Class I is defined as 0.059 $\pm$ 0.005, and Class II as $0.181 \pm 0.040$. HAT-P-5b has a Safronov number of 0.059 $\pm$ 0.005, while Class I is defined as 0.059 $\pm$ 0.005, and Class II as $0.181 \pm 0.040$. HAT-P-5b has a Safronov number of 0.059 $\pm$ 0.005, while Class I is defined as 0.059 $\pm$ 0.005, and Class II as $0.181 \pm 0.040$. HAT-P-5b has a Safronov number of 0.059 $\pm$ 0.005, while Class I is defined as 0.059 $\pm$ 0.005, and Class II as $0.181 \pm 0.040$.

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### Table 3: Orbital Fit and Planetary Parameters for the HAT-P-5b System

| Parameter       | Value |
|-----------------|-------|
| **Ephemerides:**|       |
| Period$^a$ (days) | 2.788491 $\pm$ 0.000025 |
| Transit duration (days) | 0.1217 $\pm$ 0.0012 |
| Ingress duration (days) | 0.0145 $\pm$ 0.0007 |

| RV orbital fit: |       |
| Stellar jitter$^b$ (m s$^{-1}$) | 14.4 |
| $\gamma$ (m s$^{-1}$) | 7613.8 $\pm$ 9.1 |
| $K$ (m s$^{-1}$) | 138 $\pm$ 14 |
| $e$ | 0 |

| Geometry and planetary parameters: |       |
| $a/R_*$ | 7.50 $\pm$ 0.19 |
| $R_p/R_*$ | 0.1106 $\pm$ 0.0006 |
| $a$ (AU) | 0.04075 $\pm$ 0.00076 |
| $b/R_*$ | 0.425 $\pm$ 0.048 |
| $i$ (deg) | 86.75 $\pm$ 0.44 |
| $M_p^c$ (M$_J$) | 1.06 $\pm$ 0.11 |
| $R_p^c$ ($R_J$) | 1.26 $\pm$ 0.05 |
| $\rho_p$ (g cm$^{-3}$) | 0.66 $\pm$ 0.11 |
| $g_p^e$ (m s$^{-2}$) | 16.5 $\pm$ 1.9 |

| Transit timing parameters: |       |
| $\Delta T_{\text{TR}} N_p^e$ | 6 $\pm$ 27 |
| $\Delta T_{\text{TR}} N_p^e$ | 79 $\pm$ 58 |
| $\Delta T_{\text{TR}} N_p^e$ | 6 $\pm$ 62 |
| $\Delta T_{\text{TR}} N_p^e$ | 6 $\pm$ 84 |
| $\Delta T_{\text{TR}} N_p^e$ | 11 $\pm$ 57 |

$^a$ Fixed in the orbital fit.
$^b$ Adopted (see text).
$^c$ The $\gamma$ velocity is in an absolute reference frame.
$^d$ Based on only directly observable quantities; see Southworth et al. (2007).