Title
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Permalink
https://escholarship.org/uc/item/4098w4sf

Journal
Physical review letters, 109(8)

ISSN
0031-9007

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Publication Date
2012-08-21

DOI
10.1103/physrevlett.109.081801

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Peer reviewed
Search for $tb$ Resonances in Proton-Proton Collisions at $\sqrt{s} = 7$ TeV with the ATLAS Detector

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(Received 4 May 2012; published 21 August 2012)

This Letter presents a search for $tb$ resonances in 1.04 fb$^{-1}$ of LHC proton-proton collision data collected by the ATLAS detector at a center-of-mass energy of 7 TeV. Events with a lepton, missing transverse momentum, and two jets are selected and the invariant mass of the corresponding final state is reconstructed. The search exploits the shape of the $tb$ invariant mass distribution compared to the expected standard model backgrounds. The model of a right-handed $W'$ with standard model-like couplings is chosen as the benchmark model for this search. No statistically significant excess of events is observed in the data, and upper limits on the cross section times the branching ratio of $W'_R$ resonances at 95% C.L. lie in the range of 6.1–1.0 pb for $m_{W'_R}$ masses ranging from 0.5 to 2.0 TeV. These limits are translated into a lower bound on the allowed right-handed $W'_R$ mass, giving $m_{W'_R} > 1.13$ TeV at 95% C.L.

DOI: 10.1103/PhysRevLett.109.081801

PACS numbers: 14.70.Pw, 12.60.Cn, 13.85.Rm, 14.65.Ha

This Letter presents a search for $tb$ ($t\bar{b}$ or $\bar{t}b$) resonances using data collected in 2011 by the ATLAS detector [1] at the Large Hadron Collider (LHC), corresponding to an integrated luminosity of 1.04 ± 0.04 fb$^{-1}$ [2,3] from $pp$ collisions at a center-of-mass energy of 7 TeV. These resonances include new heavy gauge bosons such as the $W'$ boson. The $W'$ boson is a charged heavy gauge boson that is predicted in many extensions of the standard model (SM) such as universal extra dimensions [4] and little Higgs models [5]. If the $W'$ boson is assumed to have similar coupling strengths to those of the SM $W$ boson, searches in the $W' \rightarrow \ell\nu$ decay channel, where $\ell$ is a charged lepton, are the most sensitive. However, the $W' \rightarrow tb$ channel is competitive if $W' \rightarrow \ell\nu$ decay is suppressed. For example, for a right-handed $W'_R$ this can happen if the right-handed neutrino $\nu_R$ is heavy enough to prevent $W'_R \rightarrow \ell\nu_R$ decay [6]. The model of a right-handed $W'_R$ with SM-like couplings is chosen as the benchmark model for the analysis presented in this Letter. The $W'_R \rightarrow tb$ decay channel has been searched for at the Tevatron [7,8]. The best previous limit on a $W'_R$ with standard model-like couplings of the $W'$ to quarks was set by the D0 experiment and excludes a $W'_R$ mass below 890 GeV at 95% confidence level.

The innermost part of the ATLAS detector [9], a tracking system in a 2 T axial magnetic field, measures the momentum of the charged particles produced in the collisions. Outside of the solenoid are the calorimeter subsystems, which measure the electron, photon, and hadronic particle energies, and the muon spectrometer, which is used to identify and measure the momentum of muons in a toroidal magnetic field. A three-level trigger system [10] reduces the event rate and selects the events for analysis.

The $tb$ resonances are searched for in the $tb \rightarrow \ell\nu bb$ decay channel, where the lepton $\ell$ is either an electron or a muon. $W'_R$ signal events are simulated to leading order (LO) with the PYTHIA v.6.421[11] Monte Carlo (MC) generator, using the MRST2007 LO* parton distribution functions (PDFs) [12]. Seven signal samples are simulated, with different $W'_R$ mass assumptions, ranging from 500 GeV to 2.0 TeV, as reported in Table I. The respective signal cross section times the branching ratio values are computed at next-to-leading-order (NLO) [13], using CTEQ6.6 PDFs [14].

Data-driven methods and MC simulated samples are used to estimate and model backgrounds. The $t\bar{t}$ process is simulated with the MC@NLO v3.4.15[15,16] MC generator, assuming a top quark mass of 172.5 GeV, and using the CTEQ6.6 PDFs. The parton shower is added using the HERWIG [17] and JIMMY [18] MC generators. The $t\bar{t}$ cross section

| $m_{W'_R}$ [GeV] | $\mathcal{B}(W'_R \rightarrow tb)$ | $\sigma \times \mathcal{B}$ [pb] |
|------------------|-------------------------------|-------------------------------|
| 500              | 0.298 ± 0.002                 | 54.6 ± 2.1                    |
| 750              | 0.319 ± 0.001                 | 10.9 ± 0.6                    |
| 1000             | 0.326 ± 0.001                 | 2.92 ± 0.18                   |
| 1250             | 0.328 < 0.001                 | 0.91 ± 0.07                   |
| 1500             | 0.330 < 0.001                 | 0.31 ± 0.03                   |
| 1750             | 0.331 < 0.001                 | 0.11 ± 0.01                   |
| 2000             | 0.332 < 0.001                 | 0.044 ± 0.005                 |

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section is obtained from the approximate next-to-next-to-leading order (NNLO) prediction calculated with the HATHOR program [19] using the MSTW2008 NNLO PDF sets [20]. The single top quark processes are simulated using the ACRMC V3.7 [21] MC generator and hadronization is performed with the PYTHIA MC generator; the cross section is calculated to approximate NNLO [22–24] using the CTEQ6.6 PDFs. Diboson processes are simulated using the HERWIG V6.5 MC generator and their cross sections are obtained at NLO using the MCFM [25] program with the MSTW2008 PDFs. The MC samples simulated with the ACRMC and HERWIG MC generators use the MRST2007 LO* PDFs. Vector boson production in association with jets (W + light jets, Wb¯b, Wc¯c, Wc, and Z + jets with up to five additional partons) is simulated using the ALPGEN V2.13 [26] MC generator, coupled with the CTEQ6L1 PDFs [14] and hadronization is performed with the HERWIG and JIMMY MC generators. In these samples, additional jets can be created from the parton shower. In order to avoid double counting between the inclusive W + n parton samples and the parton shower, overlaps are removed following the MLM matching prescription [27]. A cross section correction factor is applied to the LO W/Z + jet cross sections computed by comparing the LO and NLO predictions from the FEWZ [28] program. The Wc cross section correction factor is obtained using the MCFM [29] program with the CTEQ6.6 PDFs. All samples are passed through the full simulation of the ATLAS detector [30] based on GEANT4 [31] and are then reconstructed using the same procedure as collision data. The simulated samples include the effect of multiple pp collisions per bunch crossing (pileup) which on average is six events per bunch crossing. In order to ensure a good description of the energy scale and resolution, the trigger, the reconstruction and identification efficiency, corrections based on comparisons between data and MC events are applied to the simulated signal and background samples. The corresponding scale factors are obtained as a function of the object kinematics, resulting in final corrections of the order of a few percent.

Candidate events are identified using single high transverse momentum electron and muon triggers and stringent detector and data quality requirements. For each candidate, two jets, one isolated charged lepton, and the missing transverse momentum E_T^{miss} are required. The definition of the objects and details of a similar event selection, including lepton isolation requirements, are given in Ref. [32]. The reconstructed charged lepton is required to have a transverse momentum p_T > 25 GeV to ensure a constant trigger efficiency, |η| < 2.5 for a muon [33–35] and |η| < 2.47 for an electron [36] (the calorimeter transition region 1.37 < |η| < 1.52 is excluded), and to lie within ΔR = \sqrt{(Δη)^2 + (Δφ)^2} < 0.15 of the corresponding triggered lepton. Jets are reconstructed from energy clusters in the calorimeters with the anti-k_t algorithm [37] with a radius parameter R = 0.4 and calibrated to the hadronic energy scale [38]. Exactly two jets with p_T > 25 GeV and |η| < 2.5 are required in the event, and at least one of them must be tagged as a b jet. The b-tagging algorithm uses measurements of the impact parameters of tracks and the properties of reconstructed vertices; these are combined in a neural network to extract a tagging decision for each jet [39]. Based on a tt MC sample, the working point is chosen at a b-tagging efficiency of 57%, leading to a light-quark tagging probability of 0.2% derived from the same sample. To account for the differences between observed and simulated jet, p_T and η, distributions, the b-tagging efficiency and the corresponding scaling factors to be applied to MC simulations are derived from data [40]. Events before applying any b tagging are referred to as pretagged events. Events where one or both jets are b tagged are referred to as single- or double-tagged events, respectively.

The E_T^{miss} is calculated using calorimeter energy clusters [41] calibrated according to the reconstructed physics object to which they are associated [42]; events are required to satisfy E_T^{miss} > 25 GeV. The background contribution from multiple hadron jets (multijet background) is reduced by imposing a requirement on the sum of the W boson transverse mass m_T(W) [43] and E_T^{miss}; m_T(W) + E_T^{miss} > 60 GeV [32]. After applying all selection criteria, the acceptance times efficiency for W_h signal events with m_W = 1.0 TeV is 1.38% for single-tagged events and 0.49% for double-tagged events.

One of the most important backgrounds for the tb resonance search comes from W production in association with either heavy-flavor jets, or light-flavor jets misidentified as b jets. Multijet production is another source of background, when either a hadronic jet is misidentified as a lepton, or when a real high-p_T lepton from semileptonic decay of a heavy hadron within a jet fulfills the selection requirements. Another important background comes from tt pair production in the case that one W boson decays leptonically and the decay products of the other W boson are lost due to the detector acceptance. Other smaller backgrounds come from single top production, diboson production, and Z + jet events.

Kinematic variable distributions for the W + jet background are taken from MC samples, while the overall normalization and flavor composition are derived from data; this is done after rejecting signal-like events with the tb invariant mass m_{tb}, which is described later, satisfying m_{tb} > 500 GeV. In each jet multiplicity bin, the number of W + jet events in the data is assumed to be the difference between the number of observed data events and the number of events estimated for SM non-W + jet processes including the multijet process estimated from a data-driven method. The overall W + jet normalization factor is the ratio of the number of W + jet events in the data to the number of W + jet events in simulation. The
flavor composition of the $W + \text{jet}$ background is estimated by comparing the MC prediction to data while its dependence on jet and $b$-tagging multiplicity is modeled using MC simulations. The fractions of $Wb\bar{b}$, $Wc\bar{c}$, $Wc$, and $W + \text{light jet}$ components of the total $W + \text{jet}$ MC simulations are scaled such that the background sum equals the observed data in three separate samples: a single-tagged one-jet sample and the pretagged and single-tagged two-jet samples. The same scale factor is used for $Wb\bar{b}$ and $Wc\bar{c}$.

The multijet background normalization and the shape of each distribution are obtained from data. The shape of each multijet background distribution is taken from a data sample which requires a jet instead of an isolated lepton. This jet is required to have a detector signature similar to an electron: it must have $p_T > 25$ GeV and between 80% and 95% of its energy deposited in the electromagnetic section of the calorimeter. The jet must also be associated with at least four tracks. The normalization is estimated using a binned likelihood fit to the $E_T^{\text{miss}}$ distribution in data in which the normalization of the $W + \text{jet}$ and the multijet components is allowed to vary. The fit is performed separately in the pretagged, single-, and double-tagged samples, after applying all selection criteria except the $E_T^{\text{miss}}$ cut. The uncertainty on the multijet rate is 50% for pretagged and single-tagged events, while it amounts to 100% for double-tagged events. The uncertainty is estimated by using the $m_T(W)$ distribution instead of the $E_T^{\text{miss}}$ distribution in the binned likelihood fit, and by using multijet background models built from data samples with low and high numbers of $pp$ collisions per event.

The $t\bar{t}$, single top, $Z + \text{jet}$, and diboson events are normalized to the theoretical cross sections and the shape of each distribution is taken from the MC simulation.

Based on the theoretical predictions shown in Table I, the numbers of single- and double-tagged $W_R$ signal events expected in 1.04 fb$^{-1}$ are listed in Table II, as a function of $m_{W_R}$. Table III lists the expected background yields.

The $tb$ invariant mass is used as the observable to discriminate signal from background. The neutrino momentum in the decay $tb \rightarrow \ell \nu b\bar{b}$ is computed assuming the transverse component to be equal to $E_T^{\text{miss}}$, and extracting the longitudinal component ($p_z$) by constraining the $\ell$-$\nu$ invariant mass to $m_W = 80.42$ GeV. This gives a quadratic equation in $p_z$ and the solution with the smaller $|p_z|$ is used. If the solution is complex, only the real part is taken and the imaginary part is neglected.

Figure 1 shows the data and expected background distributions of $m_{tb}$ for single- and double-tagged two-jet events. The data event with the highest $m_{tb}$ value corresponds to a single-tagged event with $m_{tb} \approx 2.0$ TeV. The BUMPHER tool [44] is used to search for a local excess in the data due to the production of a $tb$ resonance. This tool is used to test the consistency of the data with the SM background only hypothesis, comparing the data to the SM prediction over the spectrum of the $tb$ invariant mass, scanning over sliding mass windows from 0.5 to 2.0 TeV.

The width of the mass windows is chosen to be constant in log($m_{tb}$) as shown in Fig. 1 to deal with low background MC statistics in the higher mass bins. This comparison has been performed for single- and double-tagged events separately. The region with the highest data-background difference is 1024–1129 (764–842) GeV for single (double)-tagged events. The probability of observing the SM background fluctuating up to or above the number of observed data events in these regions is 0.66 for single-tagged events and 0.72 for double-tagged events. These values, which are based on the statistical error only, indicate that there is no significant evidence for $tb$ resonances in the observed data.

Systematic uncertainties from various sources affecting the background and the signal acceptance (rate uncertainty), as well as shape changes in the invariant mass distribution (shape uncertainty) are considered.

The jet energy scale and the uncertainty on the $b$-tagging scale factors are the dominant systematic uncertainties for the signal. The background normalization yields are the dominant systematic uncertainty for the background.

| $m_{W_R}$ [GeV] | Single-tagged | Double-tagged |
|----------------|---------------|---------------|
| 500            | 973 ± 37      | 455 ± 17      |
| 750            | 174 ± 9       | 77 ± 4        |
| 1000           | 42 ± 3        | 15 ± 1        |
| 1250           | 11 ± 1        | 3.9 ± 0.3     |
| 1500           | 3.2 ± 0.3     | 1.0 ± 0.1     |
| 1750           | 1.0 ± 0.1     | 0.26 ± 0.03   |
| 2000           | 0.36 ± 0.04   | 0.09 ± 0.01   |

TABLE III. Predicted background event yields compared to the total observed event yields for single- and double-tagged two-jet events in 1.04 fb$^{-1}$ of data. All $W + \text{jet}$ samples are scaled by the factors determined from data, with the uncertainties also derived from data. The multijet estimation is from the fitting method with a 50% (100%) uncertainty for single-(double-)tagged events. All the other predictions are derived using the theoretical cross sections and uncertainties.

| Samples          | Single-tagged | Double-tagged |
|------------------|---------------|---------------|
| $W + \text{jets}$| 5970 ± 1000   | 290 ± 180     |
| Multijets        | 1120 ± 560    | 47 ± 47       |
| $t\bar{t}$       | 1560 ± 130    | 360 ± 30      |
| Single top       | 1240 ± 90     | 120 ± 10      |
| Diboson, $Z + \text{jets}$ | 320 ± 120 | 14 ± 2        |
| Total prediction | 10 220 ± 1200 | 830 ± 190     |
| Data             | 10 428        | 844           |
The distribution of $m_{tb}$ for single-tagged (top) and double-tagged (bottom) two-jet events in data compared to standard model expectations. The expected $W'$ signal, normalized to the theoretical cross section times the $b$-tagging scale factors are $p_T$ dependent and have an uncertainty between 8% and 20%. The multijet background uncertainty has already been described. The uncertainty on the normalization of the $W +$ jets background and its flavor composition include both systematic contributions and a statistical contribution from the limited size of the sample. The $W +$ jet flavor uncertainties are treated as fully correlated between $Wb\bar{b}$ and $Wc\bar{c}$ and uncorrelated otherwise. Theoretical cross section uncertainties for the top ($t\bar{t}$ and single top), diboson and $Z +$ jet backgrounds of 10%, 5%, and 60% are assigned, respectively. The $Z +$ jet theoretical cross section uncertainty is estimated based on the variation of ALPGEN parameters, and a relative uncertainty of 50% on the heavy-quark contributions, but it has a very small impact on the result due to the small contribution of $Z +$ jet events. Systematic uncertainties due to the residual differences between data and MC simulation for the reconstruction and energy calibration of jets, electrons, and muons are estimated to have a small impact on the result. The uncertainty on the integrated luminosity is 3.7% [3]. The uncertainty on the background modeling in the $m_{tb}$ distribution is evaluated using pretrained data and found to be negligible.

An uncertainty due to the MC event generator is estimated by comparing MC@NLO and POWHEG [45,46] for $t\bar{t}$ and ACERMC and MC@NLO for single top events. The uncertainty in parton shower modeling is estimated by comparing two POWHEG $t\bar{t}$ samples for which the hadronization is performed by PYTHIA or HERWIG. Uncertainties from modeling the amount of initial and final-state QCD radiation are also taken into account. The uncertainty due to the specific choice of PDFs in the simulated events is determined by reweighting the MC events using the NNPDF20, MSTW2008, and CTEQ6.6 [20] eigenvector PDF sets. Finally, an uncertainty to account for the limited MC sample sizes is also included.

No significant data excess is identified for any value of $m_{tb}$, and an upper limit on the $W_W^prime \rightarrow tb$ production cross section ($\sigma$) times the $B(W_W^prime \rightarrow tb)$ at 95% credibility-level (C.L.) is determined using a Bayesian approach assuming flat priors [47]. The likelihood function used is the product of the Poisson probabilities over all mass bins [48] per channel. The combination of single- and double-tagged events is done by extending the likelihood function; the joint likelihood is the product of Poisson probabilities over all mass bins [48] per channel. Systematic and statistical uncertainties are incorporated and treated as nuisance parameters with a Gaussian probability density function. Figure 2 shows the observed and the expected limits from single- and double-tagged events combined. Observed (expected) upper limits obtained on $\sigma(pp \rightarrow W^prime) \times B(W_W^prime \rightarrow tb)$ at 95% C.L. lie in the range...
These are currently the most stringent direct limits on \( W_R \) masses ranging from 0.5 to 2.0 TeV. These \( \sigma \times B \) limits are also applicable to a left-handed \( W \). The \( \sigma \times B \) limits are converted into mass limits using the intersection between the theoretical \( \sigma \times B \) curve as a function of \( m_{W_R} \) and the expected and observed \( \sigma \times B \) limit curves. The corresponding observed (expected) 95\% C.L. lower limit is \( m_{W_R} \geq 1.13(1.13) \) TeV. These are currently the most stringent direct limits on production of \( W_R \rightarrow tb \).

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We thank Z. Sullivan for discussions on the \( W \) model and for providing NLO signal cross section calculations. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF, DNSRC and Lundbeck Foundation, Denmark; EPLANET and ERC, European Union; IN2P3-CNRS, CEADSM/IRFU, France; GNAS, Georgia; BMBF, DFG, MPG and AvH Foundation, Germany; GSRT, Greece; ISF, MINERVA, GIF, DIP and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; RCN, Norway; MNiSW, Poland; GRICES and FCT, Portugal; MERSYS (MECTS), Romania; MES of Russia and ROSATOM, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MVZT, Slovenia; DST/NRF, South Africa; MICINN, Spain; SRC and Wallenberg Foundation, Sweden; SER, SNSF and Cantons of Bern and Geneva, Switzerland; NSC, Taiwan; TAEK, Turkey; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America. The crucial computing support from all WLCG partners is acknowledged gratefully; in particular, from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

[1] ATLAS Collaboration, JINST 3, S08003 (2008).
[2] ATLAS Collaboration, Eur. Phys. J. C 71, 1630 (2011).
[3] ATLAS Collaboration, Report No. ATLAS-CONF-2011-116, https://cdsweb.cern.ch/record/1376384.
[4] A. Datta, P.J. O’Donnell, Z. H. Lin, X. Zhang, and T. Huang, Phys. Lett. B 483, 203 (2000).
[5] M. Perelstein, Prog. Part. Nucl. Phys. 58, 247 (2007).
[6] G. Altarelli, B. Mele and M. Ruiz-Altaba, Z. Phys. C 45, 109 (1989).
[7] V.M. Abazov et al. (D0 Collaboration), Phys. Lett. B 699, 145 (2011).
[8] T. Aaltonen et al. (CDF Collaboration), Phys. Rev. Lett. 103, 041801 (2009).
[9] In the ATLAS coordinate system, the pseudorapidity \( \eta \) is defined as \( \eta = -\ln(\tan(\theta/2)) \), where \( \theta \) is measured with respect to the \( z \) axis, defined to be parallel to the beam. The azimuthal angle \( \phi \) is measured with respect to the \( x \) axis, which points toward the center of the LHC ring, and the \( y \) axis points upwards.
[10] ATLAS Collaboration, Eur. Phys. J. C 72, 1849 (2012).
[11] T. Sjostrand, S. Mrenna, and P.Z. Skands, J. High Energy Phys. 05 (2006) 026.
[12] A. Sherstnev and R.S. Thorne, Eur. Phys. J. C 55, 553 (2008).
[13] Z. Sullivan, Phys. Rev. D 66, 075011 (2002).
[14] P.M. Nadolski, H.-L. Lai, Q.-H. Cao, J. Huston, J. Pumpolin, D. Stump, W.-K. Tung, and C.-P. Yuan, Phys. Rev. D 78, 013004 (2008).
[15] S. Frixione and B.R. Webber, J. High Energy Phys. 06 (2002) 029.
[16] S. Frixione, P. Nason, and B.R. Webber, J. High Energy Phys. 08 (2003) 007.
[17] G. Corcella, I.G. Knowles, G. Marchesini, S. Moretti, K. Odagiri, P. Richardson, M.H. Seymour, and B.R. Webber, J. High Energy Phys. 01 (2001) 010.
[18] J.M. Butterworth, J.R. Forshaw, and M.H. Seymour, Z. Phys. C 72, 637 (1996).
[19] M. Aliev, H. Lackera, U. Langenfelda, S. Mochb, P. Uwera, and M. Wiedermann, Comput. Phys. Commun. 182, 1034 (2011).
[20] A.D. Martin, W.J. Stirling, R.S. Thorne, and G. Watt, Eur. Phys. J. C 63, 189 (2009).
[21] B.P. Kersevan and E. Richter-Was, arXiv:hep-ph/0405247.
[22] N. Kidonakis, Phys. Rev. D 82, 054018 (2010).
[23] N. Kidonakis, Phys. Rev. D 81, 054028 (2010).
[24] N. Kidonakis, Phys. Rev. D 83, 091503 (2011).
[25] J.M. Campbell, R.K. Ellis, and C. Williams, J. High Energy Phys. 07 (2011) 018.
[26] M.L. Mangano, M. Moretti, F. Piccinini, R. Pittau, and A.D. Polosa, J. High Energy Phys. 07 (2003) 001.
[27] S. Hoche, F. Krauss, N. Lavesson, L. Lonnblad, M. Mangano, A. Schalicke, and S. Schumann, arXiv:hep-ph/0602031v1.
[28] C. Anastasiou, L.J. Dixon, K. Melnikov, and F. Petriello, Phys. Rev. D 69, 094008 (2004).
[29] J.M. Campbell and R.K. Ellis, Phys. Rev. D 60, 113006 (1999).
[30] ATLAS Collaboration, Eur. Phys. J. C 70, 823 (2010).
[31] S. Agostinelli et al., Nucl. Instrum. Methods Phys. Res., Sect. A 506, 250 (2003).
[32] ATLAS Collaboration, Eur. Phys. J. C 71, 1577 (2011).
[33] ATLAS Collaboration, Report No. ATLAS-CONF-2011-201, https://cdsweb.cern.ch/record/1336750.
[34] ATLAS Collaboration, Report No. ATLAS-CONF-2011-046, https://cdsweb.cern.ch/record/1338575.
[35] ATLAS Collaboration, Report No. ATLAS-CONF-2011-063, https://cdsweb.cern.ch/record/1345743.
[36] ATLAS Collaboration, Eur. Phys. J. C 72, 1909 (2012).
The W boson transverse mass is defined as \( \sqrt{2p_T^2 + E_T^{miss}(1 - \cos \Delta \phi)} \), where \( p_T \) is the \( p_T \) of the lepton and \( \Delta \phi \) is the azimuthal angle separation between the lepton and \( E_T^{miss} \).
