Symmetries of the Space of Linear Symplectic Connections

Daniel J.F. FOX

Departamento de Matemáticas del Área Industrial, Escuela Técnica Superior de Ingeniería y Diseño Industrial, Universidad Politécnica de Madrid, Ronda de Valencia 3, 28012 Madrid, Spain
E-mail: daniel.fox@upm.es

Received June 30, 2016, in final form January 07, 2017; Published online January 10, 2017
https://doi.org/10.3842/SIGMA.2017.002

Abstract. There is constructed a family of Lie algebras that act in a Hamiltonian way on the symplectic affine space of linear symplectic connections on a symplectic manifold. The associated equivariant moment map is a formal sum of the Cahen–Gutt moment map, the Ricci tensor, and a translational term. The critical points of a functional constructed from it interpolate between the equations for preferred symplectic connections and the equations for critical symplectic connections. The commutative algebra of formal sums of symmetric tensors on a symplectic manifold carries a pair of compatible Poisson structures, one induced from the canonical Poisson bracket on the space of functions on the cotangent bundle polynomial in the fibers, and the other induced from the algebraic fiberwise Schouten bracket on the symmetric algebra of each fiber of the cotangent bundle. These structures are shown to be compatible, and the required Lie algebras are constructed as central extensions of their linear combinations restricted to formal sums of symmetric tensors whose first order term is a multiple of the differential of its zeroth order term.

Key words: symplectic connection; compatible Lie brackets; Hamiltonian action; symmetric tensors

2010 Mathematics Subject Classification: 53D20; 53D05; 53C05; 17B99

1 Introduction

A (linear) symplectic connection on a symplectic manifold \((M, \Omega)\) is a torsion-free affine connection preserving the symplectic form \(\Omega\). The space \(\mathcal{S}(M, \Omega)\) of symplectic connections is a symplectic affine space; the difference of two symplectic connections is identified with a section of the bundle \(S^3(T^*M)\) of completely symmetric covariant three tensors, and the symplectic form \(\Omega\) on \(\mathcal{S}(M, \Omega)\) is given by integrating over \(M\) the pairing induced on such sections by \(\Omega\). The geometry of the symplectic affine space \((\mathcal{S}(M, \Omega), \Omega)\) is the focus of this note.

Efforts have been made to identify geometrically interesting classes of symplectic connections analogous to classes of connections studied in the Riemannian setting, such as the Levi-Civita connections of Einstein or constant scalar curvature metrics. Such classes are defined in terms of the zeros and critical points of functionals on \((\mathcal{S}(M, \Omega), \Omega)\) equivariant or invariant with respect to some action of some group of symplectomorphisms or Hamiltonian diffeomorphisms.

The geometry of any symplectic space is reflected in its Poisson algebra of functions. Interesting functionals on \((\mathcal{S}(M, \Omega), \Omega)\) are constructed from the curvature of \(\nabla \in \mathcal{S}(M, \Omega)\). The simplest interesting classes of symplectic connections that have been studied are those of Ricci type, the preferred symplectic connections (these include the Ricci flat symplectic connections), and the critical symplectic connections. The focus here is on preferred and critical symplectic connections.
The critical points of the integral of the trace of the square of the Ricci endomorphism are the preferred symplectic connections first studied by F. Bourgeois and M. Cahen [5, 6]. A symplectic connection is preferred if and only if there vanishes the complete symmetrization of the covariant derivative of its Ricci tensor $\text{Ric}_{ij} = R_{ij}$, that is, it solves $(\delta^s\text{Ric})_{ijk} = -\nabla_i R_{jk} = 0$. The functionals given by integrating the trace of higher even powers of the Ricci endomorphism play a role in the averaging procedure used by Fedosov in [9], and, as is explained briefly in Section 4.3, their critical points generalize preferred symplectic connections. However, because the first Pontryagin class of a symplectic manifold equals the integral of an expression quadratic in the curvature tensor of the connection, the class of preferred symplectic connections is the only interesting class of symplectic connections obtained as the critical points of a functional given by integrating expressions quadratic in the curvature tensor (this observation is due to Proposition 2.4 of [6]; see also the introduction of [10]).

In [7], M. Cahen and S. Gutt showed that the action of the group of Hamiltonian diffeomorphisms on $S(M, \Omega)$ is Hamiltonian, with a moment map denoted here by $\mathcal{K}(\nabla)$ (see (4.3) for its definition). In [10], the author began the study of the critical symplectic connections, that are defined to be the critical points of the integral $E(\nabla) = \int_M \mathcal{K}(\nabla)^2 \Omega_n$ (where $\Omega_n = \frac{1}{n!} \Omega^n$) with respect to arbitrary variations of $\nabla$. A symplectic connection $\nabla$ is critical if and only if the Hamiltonian vector field $H_{\mathcal{K}(\nabla)}$ generated by $\mathcal{K}(\nabla)$ is an infinitesimal automorphism of $\nabla$. Explicitly this means that $\mathcal{K}(\mathcal{K}(\nabla)) = 0$ where $\mathcal{K}(f)_{ijk} = (\mathcal{L}_{H_f} \nabla)_{ij} \rho_{pk}, H_f$ is the Hamiltonian vector field generated by $f \in C^\infty(M)$, and $\mathcal{L}_{H_f}$ is the Lie derivative along $H_f$. The moment map $\mathcal{K}(\nabla)$ is the sum of a multiple of the twofold divergence of the Ricci tensor of $\nabla$ and a multiple of the complete contraction of the first Pontryagin form of $\nabla$ with $\Omega \wedge \Omega$. By the main theorem of D. Tamarkin [19], a functional on $S(M, \Omega)$ given by integrating against $\Omega_n$ polynomials in the curvature tensor and its covariant derivatives contracted with the symplectic form is independent of the choice of connection if and only if its integrand is, modulo divergences, a polynomial in Pontryagin forms contracted with the symplectic form. It follows that $\mathcal{K}$ is in some sense, not made precise here, the simplest symplectomorphism equivariant map from $S(M, \Omega)$ to $C^\infty(M)$.

Suppose $(M, \Omega)$ is compact, and fix a reference symplectic connection $\nabla_0$. Here there are constructed, for each $(s, t) \in \mathbb{R} \times \mathbb{R}^X$, a Lie bracket $\langle \langle \cdot, \cdot \rangle \rangle_{s,t}$ on

$$\tilde{\mathfrak{h}} = C^\infty(M) \oplus \Gamma(S^2(T^*M)) \oplus \Gamma(S^3(T^*M)) \quad (1.1)$$

and a Hamiltonian action of the Lie algebra $\tilde{\mathfrak{h}}, \langle \langle \cdot, \cdot \rangle \rangle_{s,t}$ on $(S(M, \Omega), \mathfrak{g})$. The associated moment map is essentially a linear combination of $\mathcal{K}$ and the Ricci tensor (see (6.16) for the precise expression when $(s, t) = (1, 1)$; it requires some preliminary discussion to make sense of it). The choice of $\nabla_0$ is unimportant in the sense that the Lie algebras associated with different choices of $\nabla_0$ are isomorphic. In fact, the brackets $\langle \langle \cdot, \cdot \rangle \rangle_{s,t}$ (and corresponding Hamiltonian actions) for which neither $s$ nor $t$ is 0 are all pairwise isomorphic, so isomorphic to $\langle \langle \cdot, \cdot \rangle \rangle = \langle \langle \cdot, \cdot \rangle \rangle_{1,1}$; the reason for considering the full family is that it interpolates between the degenerate cases $(s, t) = (1, 0)$ and $(s, t) = (0, 1)$ (this is discussed further below). In (1.1), the factor $C^\infty(M)$ should be seen as the central extension of the Lie algebra $\text{ham}(M, \Omega)$ of Hamiltonian vector fields, while $\Gamma(S^2(T^*M))$ should be seen as the Lie algebra of infinitesimal gauge transformations of the symplectic frame bundle. The Lie bracket $\langle \langle \cdot, \cdot \rangle \rangle$ combines and extends these actions.

The somewhat complicated definition of the bracket $\langle \langle \cdot, \cdot \rangle \rangle$ is summarized now. The precise definition is (6.14) of Section 6.1; see also Theorem 6.4. The algebra $\tilde{\mathfrak{g}}(M)$ whose elements are formal linear sums of completely symmetric tensors on $(M, \Omega)$ (see Section 3 for the precise meaning) carries two Lie brackets, a differential Lie bracket $\langle \langle \cdot, \cdot \rangle \rangle$ induced from the Schouten bracket on symmetric contravariant tensors on $M$ and an algebraic Lie bracket $\langle \langle \cdot, \cdot \rangle \rangle$ induced fiberwise by regarding the symmetric tensor algebra as the associated graded algebra of the Weyl algebra. The definition of $\langle \langle \cdot, \cdot \rangle \rangle$ requires Lemma 3.5, showing that these Lie brackets are compatible in the sense of bi-Hamiltonian systems; precisely, the differential, in the Lie algebra
cohomology of $(\cdot, \cdot)$, of the symmetrized covariant derivative along $\nabla$ equals the Schouten bracket $[\cdot, \cdot]$ (see (2) of Lemma 3.5). (Although it seems that this claim about the Schouten bracket must be known in some form to experts, I do not know a reference.) This means that any linear combination $[\cdot, \cdot]_{s,t} = s(\cdot, \cdot) + t[\cdot, \cdot]$ is again a Lie bracket.

The space $\hat{S}(M)$ is too big because its 1-graded part comprises all one-forms on $M$, rather than the closed one-forms, that are dual to symplectic vector fields. Because the $[\cdot, \cdot] = [\cdot, \cdot]_{1,1}$ bracket of closed one forms in $\hat{S}(M)$ need not be closed, it is inadequate to restrict to the subspace of $\hat{S}(M)$ the 1-graded part of which comprises closed one-forms (although something can be done in this direction; see Lemma 3.11). However, the subspace $\hat{H}(M)$, comprising formal sums $\alpha_0 + \alpha_1 + \cdots$ such that $\alpha_k \in \Gamma(S^k(T^*M))$ and $\alpha_1 = -d\alpha_0$ is a subalgebra of $(\hat{S}(M), [\cdot, \cdot])$, and this subalgebra plays a prominent role in the discussion. Lemmas 3.7 and 3.11 describe its structure.

Although $(\hat{S}(M), [\cdot, \cdot])$ cannot be made to act on $(\mathcal{S}(M, \Omega), \mathcal{K})$ in any obvious way, the subalgebra $\hat{H}(M)$ admits a symplectic affine action on $(\mathcal{S}(M, \Omega), \mathcal{K})$ such that the ideal $\hat{H}_4(M) = \prod_{k \geq 4} \Gamma(S^k(T^*M))$ acts trivially. This action is weakly Hamiltonian, but the cocycle obstructing equivariance can be computed explicitly. This yields a central extension that acts in a Hamiltonian manner on $(\mathcal{S}(M, \Omega), \mathcal{K})$, with moment map $\hat{h}$. Since the Lie ideal $\hat{H}_4(M)$ acts trivially, quotienting by it yields an action of the Lie algebra $(\hat{h}, ((\cdot, \cdot)))$.

The action of the group of gauge transformations of the symplectic frame bundle of $M$ on the space of affine connections, possibly having torsion, preserving $\Omega$, is Hamiltonian with a moment map constructed in the manner of Atiyah and Bott in [1]. See Theorem 5.1 for the precise statement. Using the symplectic form the Lie algebra of infinitesimal gauge transformations can be identified with the space $\Gamma(S^2(T^*M))$ equipped with a constant multiple of the algebraic Lie bracket mentioned above. The obstruction to this Lie algebra acting on $\mathcal{S}(M, \Omega)$ is given by the Schouten bracket. However, an extended bracket on $\Gamma(S^2(T^*M)) \oplus \Gamma(S^3(T^*M))$ built from the Poisson bracket and the algebraic bracket does act on $\mathcal{S}(M, \Omega)$, by symplectic affine transformations. This extended action can be combined in a coherent way with the action of $\text{ham}(M, \Omega)$ on $\mathcal{S}(M, \Omega)$. While the resulting action is only weakly Hamiltonian, the cocycle measuring nonequivariance of the moment map can be computed explicitly and from it there is constructed the Hamiltonian action of $(\hat{h}, ((\cdot, \cdot))_{s,t})$. This action combines, at the infinitesimal level, the actions of $\text{Ham}(M, \Omega)$ and the group of gauge transformations to yield a Hamiltonian action with a moment map $\hat{\mathcal{M}}$ that simultaneously extends the moment map $\mathcal{K}$ for the action of the Lie algebra $\text{ham}(M, \Omega)$ of Hamiltonian vector fields and the Atiyah–Bott moment map.

When $M$ is compact, there is constructed (see (6.17) for the precise definition) from $\hat{h}$ a functional $\mathcal{N}: \mathcal{S}(M, \Omega) \to \mathbb{R}$. There is a graded symmetric bilinear pairing on $\hat{H}(M)$ that descends to $\hat{h}$. Pairing the moment map $\hat{\mathcal{M}}$ with itself with respect to this pairing yields $\mathcal{N}$. The critical points of $\mathcal{N}$ are the solutions of

$$t^2 \mathcal{H}(\mathcal{K}(\nabla)) + s^2 \delta^* \text{Ric} = 0. \tag{1.2}$$

The equations (1.2) interpolate between the equations for preferred and critical symplectic connections. The case $(s,t) = (1,0)$ yields the equations for the preferred symplectic connections introduced by Bourgeois and Cahen [6], and the case $(s,t) = (0,1)$ recovers the equations for the critical symplectic connections introduced by the author in [10].

**Remark 1.1.** In [10] it is shown that on a 2-manifold a preferred symplectic connection is critical. Consequently, on a 2-dimensional symplectic manifold a preferred symplectic connection solves (1.2) for any $(s,t) \in \mathbb{R}^2$.

Section 3 gives the background needed to formulate and prove Lemma 3.5. Section 4 recalls the definitions and basic characterizations of critical and preferred symplectic connections. Section 5 discusses the failure of the symplectomorphism group to act in a Hamiltonian manner on $\mathcal{S}(M, \Omega)$. Finally, the construction of $((\cdot, \cdot))_{s,t}$ and its action on $\mathcal{S}(M, \Omega)$ is given in Section 6.
2 Preliminaries

2.1. Smooth means $C^\infty$, and all manifolds, bundles, sections, etc. are smooth, unless otherwise indicated. Throughout, $(M, \Omega)$ is a connected 2n-dimensional symplectic manifold oriented by the Liouville volume form $\Omega_n = \frac{1}{m!} \Omega^n$. For simplicity, $M$ is often assumed compact, although most claims generalize straightforwardly to the noncompact setting.

For a finite-dimensional real vector space $V$, $S^k(V)$ and $\Lambda^k(V)$ denote the $k$th symmetric and antisymmetric powers of $V$. Purely algebraic constructions on vector spaces extend straightforwardly to vector bundles and their spaces of sections, and the same notations will be used in both contexts. For a smooth vector bundle $E \to M$, $\Gamma(E)$ denotes the vector space of smooth sections of $E$, and $\Lambda^k(E)$ and $S^k(E)$ denote the $k$th antisymmetric and symmetric powers of $E$. By definition $\Lambda^0(E) = S^0(E) = C^\infty(M)$.

Tensors are usually indicated using the abstract index conventions (see [17] or [20]). Indices are labels and do not indicate a choice of frame. An index is in either up or down position. Up (down) indices label contravariant (covariant) tensors. For example, $\nabla_i \nabla_j X^k$ (which could also be written $(\nabla^2 X)_{ij}{}^k$) indicates the second covariant derivative of the vector field $X^k$. Were a frame $E_i$ chosen, $\nabla_i \nabla_j X^k$ would correspond to $(\nabla^2 X)(E_i, E_j) = \nabla_{E_i} \nabla_{E_j} X - \nabla_{\nabla_{E_i} E_j} X$, and not to $\nabla_{E_i} \nabla_{E_j} X$. The up index $k$ on $X^k$ is simply a label that indicates that $X$ is a contravariant vector field. The summation convention is used in the following form: repetition of a label in up and down position indicates the trace pairing. Enclosure of indices in square brackets (parentheses) indicates complete antisymmetrization (symmetrization) over the enclosed indices, so that, for example, $a^{ij} = a^{(ij)} + a^{[ij]}$ indicates the decomposition of a contravariant two-tensor into its symmetric and skew-symmetric parts. An index included between vertical bars $| |$ is omitted from an indicated (anti)symmetrization; for example $2a_{[ij]kl} = a_{ij kl} - a_{kl ij}$.

The conventions are illustrated by the following definitions. The curvature $R_{ij}{}^{k|l}$ and torsion $\tau_{ij}{}^k$ of an affine connection $\nabla$ are defined by $2 \nabla_i \nabla_j X^k = R_{ij}{}^k{}^{l|} X^l - \tau_{ij}{}^p \nabla_p X^k$. The Ricci curvature of $\nabla$ is $R_{ij} = R_{ij}{}^{p|} X^p$. (Sometimes, for readability, $\text{Ric}$ is written instead of $R_{ij}$.)

The antisymmetric bivector $\Omega^i$ inverse to $\Omega_{ij}$ is defined by $\Omega^i \Omega^j = -\delta_j^i$. Indices are raised and lowered using $\Omega_{ij}$ and $\Omega^i$ by contracting with these tensors consistently with the conventions $X_i = X^p \Omega_{pi}$ and $X^i = \Omega_{pj} X^p$. When it is raised or lowered, an index’s horizontal position is maintained. The symplectic sharp and flat operators on a vector field $X^i$ and a one-form $\alpha_i$ are defined by $X^\flat = X^p \Omega_{pi}$ and $\alpha^\sharp = \Omega^p \alpha_p$. These are inverses, so that $(X^\flat)^\sharp = X$. For a vector field $X$, $X_i$ and $X^i$ are synonyms and both notations will be used.

2.2. Let $\text{Symp}(M, \Omega)$ be the group of compactly supported symplectomorphisms of $(M, \Omega)$. Its Lie algebra $\text{symp}(M, \Omega)$ comprises compactly supported vector fields $X$ that are locally Hamiltonian, meaning $\mathfrak{L}_X \Omega = 0$; equivalently $X^\flat = \Omega(X, \cdot)$ is closed. Define the Hamiltonian vector field $H_f$ of $f \in C^\infty(M)$ by $H_f^i = -df^i = \Omega^p df_p$. Since $H_{f,g} = [H_f, H_g]$ for the Poisson bracket $\{f, g\} = \Omega(H_f, H_g)$ of $f, g \in C^\infty(M)$, the compactly supported Hamiltonian vector fields constitute a subalgebra $\mathfrak{ham}(M, \Omega) \subset \text{symp}(M, \Omega)$. For a 2k-form $\sigma$ and $\Omega_k = \frac{1}{k!} \Omega^k$,

$$\sigma \wedge \Omega_{n-k} = \frac{1}{k!2^k \Omega}(\sigma^{i_1 i_2 \cdots i_{k-1} i_k} \Omega^{i_{k-2} \cdots i_2 i_1} \Omega_{i_1 \cdots i_k}) = \frac{1}{k!2^k \Omega}(\sigma^{i_1 \cdots i_k} i_1 \cdots i_k \Omega_{n}).$$  \hfill (2.1)

By (2.1), $\{f, g\} \Omega_n = df \wedge dg \wedge \Omega_{n-1} = d(\Omega_n \wedge \Omega_{n-1})$. Since, if $M$ is compact, $\int_M \{f, g\} \Omega_n = 0$ for $f, g \in C^\infty(M)$, the subspace $C^\infty_c(M) \subset C^\infty(M)$ comprising mean zero functions is a Lie subalgebra of $C^\infty(M)$, isomorphic to $\text{ham}(M, \Omega)$ via the map $f \to H_f$. If $M$ is noncompact, then $\text{ham}(M, \Omega)$ is identified with the Lie algebra $(C^\infty_c(M), \{\cdot, \cdot\})$ of compactly supported smooth functions. The subgroup $\text{Ham}(M, \Omega) \subset \text{Symp}(M, \Omega)_0$ of Hamiltonian diffeomorphisms of $(M, \Omega)$ comprises symplectomorphisms of $(M, \Omega)$ that can be realized as the time one flow of a normalized time-dependent Hamiltonian on $M \times [-1, 1]$, where normalized means mean-zero.
if $M$ is compact and compactly supported if $M$ is noncompact. The infinitesimal generator of a flow by Hamiltonian diffeomorphisms is a Hamiltonian vector field, so that the Lie algebra of $\text{Ham}(M, \Omega)$ is $\text{ham}(M, \Omega)$.

2.3. An action of a Lie group $G$ on a symplectic manifold $(M, \Omega)$ is symplectic if $G$ acts by symplectic diffeomorphisms. The associated Lie algebra homomorphism from the Lie algebra $\mathfrak{g}$ of $G$ to $\text{symp}(M, \Omega)$, defined by $x \rightarrow X^x_p = \frac{d}{dt}|_{t=0} \exp(-tx) \cdot p$, is weakly Hamiltonian if there is a map $\mu: M \rightarrow g^*$ such that for each $x \in g$, the Hamiltonian vector field $H_{\mu(x)}$ equals the vector field $X^x$ determined by the action. If $\mu$ is equivariant with respect to the given action of $G$ on $M$ and the coadjoint action of $G$ on $g^*$, then the action is Hamiltonian and $\mu$ is a moment map. When $G$ is simply connected, this equivariance is equivalent to the requirement that when viewed as a map from $g$ to $C^\infty(M)$, the map $\mu$ be a Lie algebra homomorphism, $\{\mu(x), \mu(y)\} = \mu([x, y])$, and the induced action of $g$ is said to be Hamiltonian if it satisfies this last condition. For a weakly Hamiltonian action that is not necessarily Hamiltonian, the associated $\mu$ is called a nonequivalent moment map. Since some authors do not require equivariance in the definition of a moment map, the redundant expression equivariant moment map is sometimes used for clarity. These definitions are applied in the infinite-dimensional context, in which case they are to be understood formally.

2.4. This subsection fixes terminology related to affine actions of a Lie algebra. This is needed mainly in Section 6.1. A map $f: A \rightarrow B$ between affine spaces $A$ and $B$ is affine if there is a linear map $L(f): V \rightarrow W$, between the vector spaces $V$ and $W$ of translations of $A$ and $B$, such that $f(q)−f(p) = L(f)(q−p)$ for all $p, q \in A$. Under composition the bijective affine maps of $A$ to itself form the group $\mathbb{A}ff(A)$ of affine automorphisms of $A$, and $L: \mathbb{A}ff(A) \rightarrow \text{GL}(V)$ is a surjective homomorphism. A one-parameter subgroup through the identity in $\mathbb{A}ff(A)$ has the form $\exp(t\phi) + O(t^2)$ for some affine map $\phi: A \rightarrow V$. Consequently, the Lie algebra $\mathfrak{aff}(A)$ of $\mathbb{A}ff(A)$ is the vector space of affine maps from $A$ to $V$ equipped with the bracket $[f, g] = L(f) \circ g − L(g) \circ f$.

An affine representation of a Lie algebra $(\mathfrak{g}, [, , ])$ on the affine space $A$ is a Lie algebra homomorphism $\rho: \mathfrak{g} \rightarrow \mathfrak{aff}(A)$, that is, a linear map satisfying $\rho([a, b]) = L(\rho(a))\rho(b) − L(\rho(b))\rho(a)$ for all $a, b \in \mathfrak{g}$. An affine action of a Lie algebra $(\mathfrak{g}, [, , ])$ on the affine space $A$ is a biaffine map $\pi: \mathfrak{g} \times A \rightarrow A$ such that $\pi(0, p) = p$ for all $p \in A$ and $a \cdot (b \cdot p) − b \cdot (a \cdot p) = [a, b] \cdot p − p$ for all $p \in A$ and all $a, b \in \mathfrak{g}$. That this identity is an equality of vectors in $V$ explains the need for the $p$ on its right-hand side. That $\pi$ be biaffine means that the maps $p \rightarrow a \cdot p = \pi(a, p)$ and $a \rightarrow a \cdot p = \pi(a, p)$ are affine maps from $A$ to $A$ for all $a \in \mathfrak{g}$ and from $\mathfrak{g}$ to $A$ for all $p \in A$. (Note that a biaffine map need not be affine with respect to the product affine structure on $\mathfrak{g} \times A$.)

**Lemma 2.1.** For an affine space $A$ and a Lie algebra $\mathfrak{g}$, a map $\rho: \mathfrak{g} \rightarrow \mathfrak{aff}(A)$ is an affine representation if and only if $\pi: \mathfrak{g} \times A \rightarrow A$ defined by $\pi(a, p) = p + \rho(a)p$ is an affine action.

First one shows that $\rho$ is linear if and only if $\pi$ is biaffine and $\pi(0, p) = 0$ for all $p \in A$. Then one checks directly that, in this case, $((a \cdot (b \cdot p) − b \cdot (a \cdot p)) − ([a, b] \cdot p − p) = L(\rho(a))\rho(b) − L(\rho(b))\rho(a) − L(\rho([a, b]))p$. The representation $\rho$ is said to be associated with the affine action $\pi$. The stabilizer of $p_0 \in A$ under the affine action of $\mathfrak{g}$ means the subalgebra $\{a \in \mathfrak{g}: a \cdot p_0 = p_0\} = \{a \in \mathfrak{g}: \rho(a)p_0 = 0\}$.

A typical example is the affine action $X \cdot \nabla = \nabla + \nabla X \cdot \nabla$ of the Lie algebra of vector fields on $M$ on the space of affine connections on $M$. The associated affine representation is $\rho(X) = \nabla X \cdot \nabla$.

A symplectic affine space means an affine space $A$ equipped with a symplectic form $\Omega$ invariant under the action of the group of translations of $A$. An affine representation $\rho: \mathfrak{g} \rightarrow \mathfrak{aff}(A)$ on a symplectic affine space $(A, \Omega)$ is symplectic affine if each $\rho(a)$ is symplectic, meaning the linear part $L(\rho(a))$ satisfies $\Omega(L(\rho(a))u, v) + \Omega(u, L(\rho(a))v) = 0$ for all $a \in \mathfrak{g}$ and $u, v \in V$. 

Symmetries of the Space of Linear Symplectic Connections 5
2.5. An affine connection is \textit{symplectic} if $\nabla_\omega \Omega_{jk} = 0$. For background about symplectic connections see \cite{2} or \cite{8} and the references therein. The space $T(M, \Omega)$ of symplectic affine connections on $(M, \Omega)$ is an affine space modeled on the vector space $\Gamma(T^*M \otimes S^2(T^*M))$ of covariant tensors $\Pi_{ijk} = \Pi_{i(jk)}$ symmetric in the last two entries. Precisely, if $\nabla \in T(M, \Omega)$ and $\nabla = \nabla + \Pi_{ij}^k$ then $\nabla_\omega \Omega_{jk} = -2\Pi_{i[jk]}$, so that $\nabla \in T(M, \Omega)$ if and only if $\Pi_{i[jk]} = 0$. Note that $\Pi_{ip}^p = 0$, reflecting that a symplectic affine connection preserves the associated volume form $\Omega_n$. The affine subspace $S(M, \Omega) \subset T(M, \Omega)$ comprising torsion-free symplectic affine connections is modeled on the vector space $S^2(M)$, for the difference tensor of torsion-free connections satisfies $\Pi_{i[jk]}^k = 0$, and with $\Pi_{i[jk]} = 0$, this implies $\Pi_{ijk} = \Pi_{i(jk)}$. The space $S(M, \Omega)$ is nonempty, for if $\nabla$ is any torsion-free affine connection then $\nabla = \nabla + \frac{1}{2} \Omega_{pq} \nabla(I(i\Omega)_p)$ is symplectic. It is convenient to write $T_\nabla T(M, \Omega) = \Gamma(T^*M \otimes S^2(T^*M))$ or $T_\nabla S(M, \Omega) = S^2(M)$ for the tangent space to the affine spaces $T(M, \Omega)$ or $S(M, \Omega)$ at $\nabla$. The focus here is on $S(M, \Omega)$, and it is mainly in Section 5.1 that reference will be made to $T(M, \Omega)$. Throughout this paper \textit{symplectic connection} means a torsion-free symplectic connection, while \textit{symplectic affine connection} means a connection preserving $\Omega$ but possibly having torsion.

Since, for $\nabla \in S(M, \Omega)$, $\Omega_n$ is $\nabla$-parallel, the curvature $-R_{ij}^p$ of the connection induced on $\Lambda^{2n}(T^*M)$ by $\nabla$ vanishes. Hence $\nabla$ has symmetric Ricci tensor, for $2R_{ij} = -R_{ij}^p = 0$. By the Ricci identity, $0 = 2\nabla_i \nabla_j \Omega_{kl} = -2R_{ijkl}$, where $R_{ijkl} = R_{ij}^k \Omega_{kl}$. The algebraic Bianchi identity and $R_{ijkl} = 0$ yield $R_{ij}^p = -2R_{ij}^p = 2R_{ij}$, from which it follows that every nontrivial trace of $R_{ijkl}$ is a constant multiple of $R_{ij}$.

2.6. Let $(\mathcal{V}, \Omega)$ be a symplectic vector space. On $W^{p,q}(\mathcal{V}) = \Lambda^p(\mathcal{V}) \otimes S^q(\mathcal{V})$ define

$$\Omega(\alpha, \beta) = (-1)^p \frac{1}{p!} \alpha_{i_1 \ldots i_p} \beta^{j_1 \ldots j_q} \Omega_{i_1 \ldots i_p j_1 \ldots j_q}.$$ \hspace{1cm} (2.2)

This pairing is graded symmetric in the sense that $\Omega(\alpha, \beta) = (-1)^{|\alpha||\beta|} \Omega(\beta, \alpha)$ for homogeneous elements $\alpha$ and $\beta$, where $|\alpha| = p + q$ is the degree of $\alpha \in W^{p,q}(\mathcal{V})$. Integrating (2.2) yields a Symp $(M, \Omega)$-invariant graded symmetric pairing on $\Gamma(W^{p,q}(T^*M))$ defined by

$$\langle \alpha, \beta \rangle = \int_M \Omega(\alpha, \beta) \Omega_n = (-1)^{|\alpha|} \frac{\Omega_n}{p!} \int_M \alpha_{i_1 \ldots i_p} \beta^{j_1 \ldots j_q} \Omega_{i_1 \ldots i_p j_1 \ldots j_q}.$$ \hspace{1cm} (2.3)

Using an almost complex structure compatible with $\Omega$ it is straightforward to show that the pairing $\langle \alpha, \beta \rangle$ is weakly nondegenerate in the sense that if $\langle \alpha, \beta \rangle = 0$ for all compactly supported $\beta \in \Gamma(W^{p,q}(T^*M))$ then $\alpha = 0$.

2.7. Fix a $2n$-dimensional symplectic vector space $(\mathcal{V}, \Omega)$. The (linear) \textit{symplectic frame bundle} $\pi: \mathcal{F} \to M$ is the $G = Sp(\mathcal{V}, \Omega)$ principal bundle whose fiber over $p \in M$ comprises symplectic linear maps $u: (\mathcal{V}, \Omega) \to (T_p M, \Omega)$. A symplectic affine connection $\nabla$ determines and is determined by a principal connection on $\mathcal{F}$. The space $T(M, \Omega)$ of symplectic affine connections on $M$ is an affine space modeled on the vector space of 1-forms taking values in the adjoint bundle $Ad(\mathcal{F}) = \mathcal{F} \times Ad g$ associated with $\mathcal{F}$ by the adjoint action of $G$ on its Lie algebra $g = \{ \mathcal{V}, \Omega \}$.

2.8. Given $\alpha, \beta \in T_\nabla \Omega(M, \Omega)$, composing the bundle-valued two-form $(\alpha \wedge \beta)_{ijkl} = 2\alpha_{[i[p|} \beta_{j]k]}{^p}$ $= -2\alpha_{[i[p|} \beta_{j]k]}{^p}$ with the trace on endomorphisms yields the two-form $tr(\alpha \wedge \beta)_{ij} = 2\alpha_{[i[p|} q_{j]q]}{^p}$. By (2.1), $tr(\alpha \wedge \beta) \wedge \Omega_{n-1} = -\alpha_{ijk} \beta_{ij} \Omega_n$. By (2.3) the two-form $\Omega$ defined on $T(M, \Omega)$ by

$$\Omega(\alpha, \beta) = -\int_M tr(\alpha \wedge \beta) \wedge \Omega_{n-1} = \int_M \alpha_{ijk} \beta_{ij} \Omega_n = \langle \alpha, \beta \rangle, \quad \alpha, \beta \in T_\nabla \Omega,$$
is weakly nondegenerate. The translation invariance, \( \Omega_{\nabla + \Pi} = \Omega_{\nabla} \), for \( \Pi, T_{\nabla} V(M, \Omega) \) means \( \Omega \) is parallel, so closed, and so \( \Omega \) is a symplectic form. The subspace \( S(M, \Omega) \subset T(M, \Omega) \) is symplectic.

The group of diffeomorphisms acts on the space of affine connections on the right, by pullback. Precisely \( \phi^*(V) \) is defined by \( \phi^*(V)_X Y = T_{\phi^{-1}}(\nabla_{T_{\phi}(X)} T_{\phi}(Y)) \). This action preserves the torsion-free connections. Consequently, \( \text{Symp}(M, \Omega) \) acts on \( S(M, \Omega) \) and \( T(M, \Omega) \) by pullback, and this action preserves \( \Omega \). Integration by parts yields the corresponding infinitesimal statement, \( \Omega_{\nabla}(\mathcal{L}_{X} \alpha, \beta) + \Omega_{\nabla}(\alpha, \mathcal{L}_{X} \beta) = 0 \), for \( X \in \text{symp}(M, \Omega) \).

### 3 Symmetric tensors on symplectic manifolds

#### 3.1. The symmetric tensor algebra \( S(V) \), which is, by definition, a quotient of the tensor algebra, can be identified, as a graded vector space, with the graded vector space \( \oplus_{k \geq 0} S^k(V) \) equipped with the symmetric product \( (\alpha \otimes \beta)_{i_1, \ldots, i_{k+l}} = \alpha_{i_1, \ldots, i_k} \beta_{i_{k+1}, \ldots, i_{k+l}} \), for \( \alpha \in S^k(V) \) and \( \beta \in S^l(V) \). The subspace \( S^k(V) \) is identified with the space \( \text{Pol}^k(V^*) \) of homogeneous degree \( k \) polynomials on the dual vector space \( V^* \), and the product \( \otimes \) corresponds to multiplication of polynomials. Elements of the formally completed symmetric tensor algebra \( \hat{S}(V) = \prod_{k \geq 0} S^k(V) \) are identified with formal infinite sums of symmetric tensors, and correspond with formal power series on \( V^* \).

Let \( S(T^*M) = \oplus_{k \geq 0} S^k(T^*M) \) and \( \hat{S}(T^*M) = \prod_{k \geq 0} S^k(T^*M) \). Set \( S^k(M) = \Gamma(S^k(T^*M)) \). It is more convenient to work with the direct product \( \hat{S}(M) = \prod_{k \geq 0} S^k(M) \) than with \( \Gamma(S(T^*M)) \) or the direct sum \( S(M) = \oplus_{k \geq 0} S^k(M) \). Because a section of \( S(T^*M) \) need not have bounded degree, the space \( \Gamma(S(T^*M)) \) can be larger than \( S(M) \). On the other hand, the direct product \( \hat{S}(M) \) equals the space of sections of \( \hat{S}(T^*M) \). An element of \( \hat{S}(M) \) is regarded as a formal infinite sum of symmetric tensors on \( M \). There are inclusions, \( S(M) \subset \Gamma(S(M)) \subset \hat{S}(M) \), as associative algebras. In general, the algebraic constructions considered here are valid on any of these spaces, and, although they are usually formulated in terms of \( \hat{S}(M) \), their reformulations for the smaller spaces are generally trivial or straightforward.

An element \( \alpha \in S^k(V) \) or \( \alpha \in S^k(M) \) is said to have degree \( |\alpha| = k \). For all \( i \geq 1 \), the subspaces \( S_i(V) = \oplus_{i \geq 1} S^i(V) \) and \( \hat{S}_i(M) = \prod_{i \geq 1} S^i(M) \) are ideals with respect to the commutative algebra structure determined by \( \otimes \). The \( p \) = 0 case of (2.2) extends to a graded symmetric pairing on \( \hat{S}(V) \) in such a way that \( S^k(V) \) and \( S^l(V) \) are orthogonal if \( k \neq l \). Likewise, as commutative algebras, \( S(M) \) and \( \hat{S}(M) \) are filtered by the ideals \( S_k(M) \) and \( \hat{S}_k(M) \) comprising sums of tensors of degree at least \( k \), and the pairing (2.3) extends to a graded symmetric pairing on \( \hat{S}(M) \).

#### 3.2. If \( (V, \Omega) \) is a symplectic vector space, then \( S(V^*) \) carries a Lie bracket, determined with \( \otimes \) a Poisson algebra, that can be defined most simply as the unique Poisson bracket \( \{ \cdot, \cdot \} \) on \( (S(V^*), \otimes) \) such that \( (\alpha, \beta) = \Omega^{ij} \alpha_i \beta_j \) for all \( \alpha, \beta \in V^* \). By definition, \( \{ \cdot, \cdot \} \) has degree \(-2 \). The bracket \( \{ \cdot, \cdot \} \) can be described explicitly as follows. The space of functions on \( V \) carries the Poisson structure determined by \( \Omega \). Transporting this Poisson structure to \( S(V^*) \) via the identification of \( S(V^*) \) with the ring of polynomials on \( V \), yields the bracket \( \{ \cdot, \cdot \} \) on \( S(V^*) \). Concretely, for \( \alpha \in S^k(V^*) \), the corresponding polynomial \( \hat{\alpha} \in \text{Pol}^k(V) \) is \( \hat{\alpha}(u) = \alpha_{i_1, \ldots, i_k} u^{i_1} \cdots u^{i_k} \), and it generates the Hamiltonian vector field \( H_{\hat{\alpha}} = -k \alpha^j_{i_1, \ldots, i_{k-1}} u^{i_1} \cdots u^{i_k-1} \) on \( V \). The Poisson bracket of \( \hat{\alpha} \) and \( \hat{\beta} \) is the element \( \Omega(H_{\hat{\alpha}}, H_{\hat{\beta}}) \) of \( \text{Pol}^{k+l-2}(V) \) corresponding to \( (\alpha, \beta) \in S^{k+l-2}(V^*) \) defined by

\[
(\alpha, \beta)_{i_1, \ldots, i_{k+l-2}} = \frac{k!}{k+l-2!} \left( (l-1) \alpha_{p(i_1, \ldots, i_{k-1})} \beta_{i_{k}, \ldots, i_{k+l-2}} - (l-1) \beta_{p(i_1, \ldots, i_{k+1})} \alpha_{i_{k}, \ldots, i_{k+l-2}} \right).
\]

In the case \( l = 1 \) and \( k > 1 \), \( (\alpha, \beta) = k \beta^p \alpha_{p(i, \ldots, i_{k-1})} \) is simply \( k \) times interior multiplication of the vector field \( \beta^i \) in \( \alpha \). The Poisson bracket (3.1) extends to \( \hat{S}(V^*) \) with the same definition.
Although it is true by construction that \((\cdot, \cdot)\) is a Lie bracket that forms with the symmetric product \(\otimes\) a Poisson structure, this can also be checked directly. Let \(\gamma \in S^m(V^*)\). That \((\cdot, \cdot)\) is a Lie bracket follows by summing cyclic permutations of the identity
\[
((\alpha, \beta), \gamma)_{i_1\cdots i_{k+l+m-4}} = klm ((k-1)\alpha p q_{i_1\cdots i_{k-2}} \beta_{i_1\cdots i_{k-1}} \cdots \beta_{i_{k+l} \cdots i_{k+l+m-4}} p_{i_{k+l+1} \cdots i_{k+l+m-4}} q_

\]
(obtained by using both equalities in (3.1)). Similarly, verifying \((\alpha \otimes \beta, \gamma) = \alpha \circ (\beta, \gamma) + (\alpha, \gamma) \otimes \beta\) is straightforward using (3.1), and so \((\cdot, \cdot)\) is a Poisson bracket.

Note that \((S^2(V^*), \cdot, \cdot)\) is a Lie subalgebra of \((S(V^*), \cdot, \cdot, \cdot)\). Using the easily checked invariance \(\Omega((\alpha, \beta, \gamma) + \Omega(\beta, (\alpha, \gamma)) = 0\) for \(\alpha \in S^2(V^*)\) and \(\beta, \gamma \in S^k(V^*)\), it is straightforward to show that the map \(S^2(V^*) \to \text{End}(V)\) defined by \(\alpha \to (\alpha, \cdot)|_S^1(\gamma, V)\) is a Lie algebra isomorphism onto the symplectic Lie algebra \((V^*, \Omega)\) (see Section 2 of [15] for a detailed discussion).

**Remark 3.1.** The Weyl algebra of the symplectic vector space \((V, \Omega)\) is the quotient of the tensor algebra of \(V\) by the ideal generated by elements of the form \(u \otimes v - v \otimes u - \Omega(u, v)\) for \(u, v \in V\). The additive filtration of the Weyl algebra at level \(k\) is spanned by \(k\)-fold products of elements of \(V\). Its associated graded algebra is isomorphic as a vector space to the symmetric algebra of \(V\) and carries a Poisson bracket defined by projecting the commutator of elements of level \(k\) and \(l\) onto level \(k + l - 2\). Applying this construction to \((V^*, \Omega)\) yields the algebraic bracket (3.1).

**Remark 3.2.** Another construction of (3.1) can be given in terms of the notion of prolongation. This is recalled following [12]; see also [18, Chapter 7.3] and [14, Chapter 1.1].

The first prolongation \(g^{(1)}\) of the subspace \(g \subseteq \text{End}(V, \omega) \simeq \omega \otimes V^*\) is defined by \(g^{(1)} = \{A \in \text{hom}(V, g) \simeq g \otimes V^* : A(u) v = A(v) u\} for all u, v \in V\}, \) and, for \(k \geq 2\), the \(k\)th prolongation \(g^{(k)}\) of \(g\) is defined inductively by \(g^{(k)} = (g^{(k-1)})^{(1)}\). As vector spaces,
\[
g^{(k)} = \left(g \otimes \otimes (V^*)\right) \cap (\omega \otimes S^{k+1}(V^*)).
\]
When \(\omega = V\), so \(g \subseteq \text{End}(V)\), define \(g^{(0)} = g\) and \(g^{(-1)} = V\). Using the identification (3.2), a graded Lie bracket is defined on the direct sum \(\bigoplus_{k \geq -1} g^{(k)}\) and its formal completion \(\prod_{k \geq -1} g^{(k)}\) as follows. Define an antisymmetric bilinear pairing \([\cdot, \cdot] : \omega \otimes S^{k+1}(V^*) \times \omega \otimes S^{k+1}(V^*) \to \omega \otimes S^{k+1}(V^*)\) by setting
\[
[v_1 \otimes s_1, v_2 \otimes s_2] = v_1 \otimes ((D_{v_2}s_1) \otimes s_2) - v_2 \otimes (s_1 \otimes (D_{v_1} s_2)),
\]
for \(v_1, v_2 \in \omega\), \(s_1 \in S^k(V^*)\), and \(s_2 \in S^l(V^*)\), and extending linearly, where \(D_v : S^k(V^*) \to S^{k-1}(V^*)\) is defined by \(D_v s = k t(v)s\) for \(v \in \omega\) (the constant factor is chosen so that \(D_v\) is a derivation of the symmetric product \(\otimes\)). The bracket (3.3) restricts to yield a bracket \([\cdot, \cdot] : g^{(k)} \times g^{(l)} \to \omega \otimes S^{k+l+1}(V^*)\), and it is straightforward to check that the image is contained in \(g^{(k+l)}\) and that the resulting bracket satisfies the Jacobi identity, so makes \(\bigoplus_{k \geq -1} g^{(k)}\) and \(\prod_{k \geq -1} g^{(k)}\) into Lie algebras. In the case \((V, \omega)\) is a symplectic vector space and \(g = (V, \omega)\), this can be seen as follows. If the differential of the flow of a vector field on \(V\) preserves the symplectic frame bundle \(V \times Sp(V, \Omega)\) then the \((k+1)\)st coefficient of the Taylor expansion of the vector field takes values in \(g^{(k)}\). The Lie bracket is that induced by taking jets of Lie brackets of vector fields. Since \(\omega = V\), an element of \(g^{(0)}\) can be viewed as an endomorphism of \(V\); the definition (3.3) is made so that the bracket \([\cdot, \cdot]\) agrees on \(g^{(0)}\) with the algebraic commutator of endomorphisms. That is \([A, B]_i = (A \circ B - B \circ A)_i = A_p B_{ij} p - B_p A_{ij} p\) for \(A, B \in g^{(0)} \subset \text{hom}(V, V)\). An endomorphism \(A_i\) of \((V, \omega)\) is infinitesimally symplectic if and only if \(A_{ij} = 0\), so, using \(\omega\), \(g\) is identified with \(S^2(V^*)\) equipped with the Lie bracket \([\alpha, \beta]_{ij} = 2\alpha_{p(i} \beta_{j)} p = 2(\alpha, \beta)_{ij}\) for
α, β ∈ S^2(ν^*). Likewise, using Ω to identify the leading V (= ω) in (3.2) with V^*, it follows from (3.2) that the kth prolongation g^{(k)} is identified with S^{k+2}(V^*). Then Θ_{k≥−1}g^{(k)} and \prod_{k≥−1}g^{(k)} are identified with the ideals S_1(V^*) ⊂ S(V^*) and \tilde{S}_1(V^*) ⊂ S(V^*) and (3.3) induces on these an algebraic Lie bracket [·, ·] of degree −2 with respect to the grading of S(V). The bracket [·, ·] is related to the bracket (3.1) by kl[α, β] = (k + l − 2)(α, β) for α ∈ S^k(V^*) and β ∈ S^l(V^*).

The preceding can be recast in the following manner. The graded linear map η: S(V^*) → S(V^*) defined by η(α) = kα for α ∈ S^k(V^*) restricts to a linear isomorphism on \tilde{S}_1(V^*). Then

\{0\} → R \overset{ι}{\rightarrow} (S(V^*), (·, ·)) \overset{η}{\rightarrow} (\tilde{S}_1(V^*), [·, ·]) \rightarrow \{0\}

is a central extension of Lie algebras, where ι is the inclusion of R as the central ideal S^0(V^*) ⊂ (S(V^*), (·, ·)). The map η is an algebraic version of the operator associating to the polynomial function \hat{α}(u) = α_{i_1...i_k}u^{i_1}...u^{i_k} the Hamiltonian vector field \hat{H}_α^i = −kα_{i_1...i_{k−1}}u^{i_1}...u^{i_{k−1}}. Its restriction to \tilde{S}_1(V^*) is an isomorphism between (·, ·) and [·, ·].

Remark 3.3. A parallel distribution on a symplectic fibration such that the corresponding parallel transport preserves the fiberwise symplectic structure is also called a symplectic connection (see for example [16]). The symplectic affine connections considered here are symplectic connections in this more general sense, although the converse is clearly false. For this reason (and also to honor their role in Fedosov’s deformation quantization) what are here called symplectic connections are sometimes called Fedosov connections; here the notions are distinguished terminologically as linear and nonlinear symplectic connections.

The two notions encapsulate different points of view. Each is equivalent to the data of a principal connection on what might be called the symplectic frame bundle of M, the point being that symplectic frame bundle has two possible meanings, depending on whether symplectic refers to the symplectic linear group or the symplectomorphism group of the reference symplectic vector space (V, Ω). The linear symplectic frame bundle is the reduction of the usual linear frame bundle determined by restricting to symplectic frames. A much flabbier notion is obtained by instead considering as frames at p ∈ M all symplectomorphisms from (V, Ω) to (T_pM, Ω) mapping 0 ∈ V to 0 ∈ T_pM. Although to define such a nonlinear symplectic frame bundle rigorously requires restricting the class of symplectomorphisms considered, or working formally, e.g., with infinite jets, morally there results a principal bundle with structure group an infinite-dimensional group of (perhaps formal) symplectomorphisms of (V, Ω), and, by construction, \mathcal{F} is a reduction of this bundle. Correspondingly, there are two quite distinct notions of symmetries of S(M, Ω). The usual one is to consider principal bundle automorphisms of the linear symplectic frame bundle \mathcal{F}, regarded as a principal bundle for the finite-dimensional group G = Sp(V, Ω) of linear symplectic transformations. On the other hand, the symplectic frame bundle \mathcal{F} can be regarded as a subbundle of a principal bundle with structure group some infinite-dimensional Lie group \hat{G} of symplectomorphisms of (V, Ω) fixing the origin.

The automorphism group of \mathcal{F} comprises the G-equivariant bundle maps from \mathcal{F} to itself. The group \mathcal{G}(\mathcal{F}) of gauge transformations of \mathcal{F} is its subgroup comprising bundle automorphisms of \mathcal{F} covering the identity. A gauge transformation is naturally identified with a map from \mathcal{F} to G, equivariant with respect to the action of G on itself by conjugation, or, equivalently, with a section of the bundle \mathcal{F} ×_G G associated with \mathcal{F} by this action. Via this identification the group structure corresponds to fiberwise composition. An infinitesimal gauge transformation is a section of the adjoint bundle Ad(\mathcal{F}). These form a Lie algebra \mathcal{L}(\mathcal{F}) under fiberwise Lie bracket.

In the case of a frame bundle, a (infinitesimal) gauge transformation can be identified with a section of the bundle \text{End}(TM) of endomorphisms of the tangent bundle. Precisely, an element of \mathcal{G}(\mathcal{F}) is identified with \gamma ∈ \Gamma(\text{End}(TM)) such that \gamma q g^q_j = g_i^p g_j^q \Omega_{pq} = \Omega_{ij}. The inverse
section \( g^{-1} \) then satisfies \((g^{-1})_{ij} = -g_{ji} \). Similarly, an infinitesimal gauge transformation of \( \mathcal{F} \) is identified with a section \( \alpha_i^j \in \Gamma(\text{End}(TM)) \) such that \( \alpha_{[ij]} = 0 \); in general it is more convenient to identify \( \alpha \) with the corresponding element \( \alpha_{ij} \) of \( S^2(M) \). Alternatively, the Lie algebra \( g \) is isomorphic as a \( G \)-module to \( S^2(V^*) \) via the map \( A \to A^\gamma \) defined by \( A^\gamma(x, y) = \Omega(Ax, y) \), and via this isomorphism \( \Gamma(\text{Ad}(\mathcal{F})) \) is identified with \( \langle S^2(M), [\cdot, \cdot] \rangle \).

The preceding all makes formal sense with the nonlinear symplectic frame bundle in place of \( \mathcal{F} \), and with the corresponding structure group \( \mathcal{G} \) in place of \( G \). Such notions can be given rigorous sense at the Lie algebra level. The formal analogue of the infinitesimal \( \mathcal{G} \)-gauge transformations is the Lie algebra \( \hat{S}_2(M), [\cdot, \cdot] \), corresponding to the fiberwise action of the truncated prolongation \( (\prod_{i \geq 0} g^{(i)})[\cdot, \cdot] \) (as in Remark 3.2), comprising formal sums of covariant symmetric tensors of rank at least two.

### 3.3.

The pairing \((3.3)\) on \( S^k(M) \) extends to a \( \text{Symp}(M, \Omega) \)-invariant bilinear pairing on \( \hat{S}(M) \) by declaring \( S^k(M) \) and \( S^l(M) \) orthogonal if \( k \neq l \). This pairing is graded symmetric in the sense that \( \langle \alpha, \beta \rangle = (-1)^{|\alpha||\beta|} \langle \beta, \alpha \rangle \) for homogeneous elements \( \alpha \) and \( \beta \). If a function is regarded as a 0-tensor, then \( \langle \cdot, \cdot \rangle \) agrees with the \( L^2 \) inner product \( \langle f, g \rangle = f \int_M g \Omega \) on \( C^\infty(M) \).

A linear operator \( \mathcal{P} : S^p(M) \to S^q(M) \) has degree \(|\mathcal{P}| = q - p \). More generally, a linear operator \( \mathcal{P} : S(M) \to S(M) \) has degree \(|\mathcal{P}| = p \) if \( \mathcal{P}(\alpha) \in S^{k+p}(M) \) whenever \( \alpha \in S^k(M) \). The \textit{(formal) adjoint} \( \mathcal{P}^* : S^q(M) \to S^p(M) \) of \( \mathcal{P} \) is the degree \(-|\mathcal{P}| \) operator defined by \( \langle \mathcal{P}(\alpha), \beta \rangle = (-1)^{|\alpha||\beta|} \langle \alpha, \mathcal{P}^* \beta \rangle \). The choice of sign guarantees that \( (\mathcal{P}^*)^* = \mathcal{P} \) and \( (\mathcal{P} \mathcal{Q})^* = (-1)^{|\mathcal{P}||\mathcal{Q}|} \mathcal{Q}^* \mathcal{P}^* \).

For example, for \( \alpha \in S^2(M), (\alpha, \cdot)^* = -\langle \alpha, \cdot \rangle \), because the pairing \( \langle \cdot, \cdot \rangle \) is invariant with respect to the linear map \( \langle \alpha, \cdot \rangle : S(M) \to S(M) \) in the sense that
\[
\langle (\alpha, \Pi_1), \Pi_2 \rangle + \langle \Pi_1, (\alpha, \Pi_2) \rangle = 0, \quad \text{for } \Pi_1, \Pi_2 \in S^k(M). \tag{3.4}
\]

### 3.4.

For \( \nabla \in \mathcal{S}(M, \Omega) \) and \( \alpha \in S^k(M) \) define \( d_{\nabla} \alpha_{i_1 \ldots i_{k+1}} \) by \( d_{\nabla} \alpha_{i_1 \ldots i_{k+1}} = 2 \nabla_{[i_1} \alpha_{i_2 i_3 \ldots i_{k+1}]} \), and define the degree \(-1 \) \textit{symplectic divergence} operator \( \delta_{\nabla} : \hat{S}(M) \to \hat{S}(M) \) by
\[
\delta_{\nabla} \alpha_{i_1 \ldots i_{k+1}} = (-1)^{k-1} \frac{1}{2} (d_{\nabla} \alpha)^{p \ i_1 \ldots i_{k+1}} - (-1)^k \nabla_p \alpha_{i_1 \ldots i_{k+1}} p, \quad \alpha \in S^k(M).
\]

The formal adjoint \( \delta_{\nabla}^* : S^k(M) \to S^{k+1}(M) \) of \( \delta_{\nabla} \) is given by
\[
\delta_{\nabla}^* \alpha_{i_1 \ldots i_{k+1}} = -\nabla_{[i_1} \alpha_{i_2 i_3 \ldots i_{k+1}]} - \nabla_{i_1} \alpha_{i_2 i_3 \ldots i_{k+1}} + \frac{k+1}{k-1} d_{\nabla} \alpha_{i_1 (i_2 i_3 \ldots i_{k+1})} \quad (3.5)
\]

When it is not necessary to indicate the dependence on \( \nabla \) the subscripts are omitted and there are written \( \delta \) and \( \delta^* \) instead of \( \delta_{\nabla} \) and \( \delta_{\nabla}^* \). When \( \delta^* \) is applied to functions the subscript is always omitted, because \( \delta^* f = -df = H_f \). Note that the Poisson bracket \( \{ f, g \} \) of \( f, g, \in C^\infty(M) \) is expressible as \( \{ f, g \} = (\delta^* f, \delta^* g) \).

Lemma 3.4 is needed in the proof of Lemma 5.3.

### Lemma 3.4.

For \( \nabla \in \mathcal{S}(M, \Omega) \) and \( \alpha \in S^k(M) \),
\[
(k+1) \delta^* \alpha_{\ldots i_k} + k \delta^* \delta \alpha_{\ldots i_k} = (-1)^k (2k R_{i_1 \ldots i_{k+1}} p^{\alpha_{i_2 \ldots i_{k+1}}} \alpha_{\ldots i_k}) - k(k-1) R_{i_1 \ldots i_{k+1}} p^{\alpha_{i_2 \ldots i_{k+1}} \alpha_{\ldots i_k}} \tag{3.5}
\]

\[
\text{Proof.} \quad \text{The Ricci identity yields}
\]
\[
(-1)^k \nabla_{\alpha_{i_1 \ldots i_k}} = \nabla_{i_1} \nabla p^{\alpha_{i_2 \ldots i_k}} = \nabla p^{\nabla_{i_1} \alpha_{i_2 \ldots i_k}} + (k-1) R_{i_1 \alpha_{i_2 \ldots i_k}} p^{\alpha_{i_2 \ldots i_k} \alpha_{\ldots i_k}} - R_{i_1 \alpha_{i_2 \ldots i_k}} p^{\alpha_{i_2 \ldots i_k} \alpha_{\ldots i_k}} \cdot \tag{3.6}
\]

Combining (3.6) with
\[
(-1)^k (k+1) \delta^* \alpha_{\ldots i_k} = (k+1) \nabla p^{\nabla_{i_1} \alpha_{i_2 \ldots i_k}} = k \nabla p^{\nabla_{i_1} \alpha_{i_2 \ldots i_k}} + \nabla p^{\nabla_{i_1} \alpha_{i_2 \ldots i_k}}
\]

yields (3.5).
3.5. The Schouten bracket of contravariant symmetric tensors on a smooth manifold \( M \) is defined by regarding such tensors as functions on the cotangent bundle \( T^*M \) polynomial in the fibers and pairing them via the canonical Poisson bracket. When \( M \) is symplectic the resulting Poisson algebra structure can be transferred via the symplectic form to the algebra of covariant symmetric tensors. The resulting degree \(-1\) pairing \([\cdot, \cdot] : S^k(M) \times S^l(M) \to S^{k+l-1}(M)\) is expressible in terms of any symplectic connection \( \nabla \in \mathcal{S}(M, \Omega) \) by

\[
[\alpha, \beta]_{i_1 \ldots i_{k+l-1}} = -k\alpha_{p(i_1 \ldots i_{k-1}} \nabla^p \beta_{i_{k-1} \ldots i_{k+l-1})} + l\beta_{p(i_1 \ldots i_{l-1}} \nabla^p \alpha_{i_{l-1} \ldots i_{k+l-1})}.
\]  

(3.7)

By the definition (3.7), the Schouten bracket of functions is trivial. In (3.7) the signs are chosen so that \([X^\gamma, Y^\gamma] = [X, Y]^\gamma\) for \( X, Y \in \Gamma(TM) \). More generally, if \( X \in \text{symp}(M, \Omega) \) then \([X^\gamma, \alpha] = \mathcal{L}_X \alpha\). In particular, \( \mathcal{L}_H \alpha = [\alpha, df] = [\delta^* f, \alpha] \) for \( f \in C^\infty(M) \). Since \( \text{symp}(M, \Omega) \) is a Lie algebra, the Schouten bracket of closed one-forms is a closed one-form. The operator \( \delta^* \) is a homomorphism from \((C^\infty(M), \{\cdot, \cdot\})\) to \((\text{ham}(M, \Omega), [\cdot, \cdot])\), as

\[
[\delta^* f, \delta^* g] = [df, dg] = -d\{f, g\} = \delta^* \{f, g\}, \quad \text{for } f, g \in C^\infty(M).
\]  

(3.8)

Since the pairing (3.7) does not depend on the choice of \( \nabla \in \mathcal{S}(M, \Omega) \), it should be regarded as an object of a differential topological character attached to the symplectic manifold. In particular, the Schouten bracket is \( \text{Symp}(M, \Omega)\)-equivariant. The corresponding infinitesimal statement, the infinitesimal symplectomorphism equivariance of the Schouten bracket, is equivalent to the statement that the Lie derivative \( \mathcal{L}_X \) along \( X \in \text{symp}(M, \Omega) \) is a derivation of \([\cdot, \cdot]\). Alternatively this is a consequence of the identity \([X^\gamma, [\alpha, \beta]] = \mathcal{L}_X [\alpha, \beta]\) in conjunction with the Jacobi identity.

3.6. Two Lie brackets \([\cdot, \cdot]\) and \([\cdot, \cdot]\) on a vector space \( g \) are **compatible** if any linear combination of them is again a Lie bracket. For background on compatible Lie brackets see, for example, [3, 4, 11], and the references therein. A straightforward computation shows that compatibility of two Lie brackets is equivalent to the condition that each of the brackets is a cocycle with respect to the other; this means \([\cdot, \cdot]\) is a cocycle of the cochain complex of \((g, [\cdot, \cdot])\) with coefficients in its adjoint representation, and likewise with \([\cdot, \cdot]\) and \([\cdot, \cdot]\) interchanged. From this second characterization it follows that if one of two Lie brackets is a coboundary in the Lie algebra cohomology of the other, then the brackets are compatible.

**Lemma 3.5.** Let \((M, \Omega)\) be a symplectic manifold. For \( \nabla \in \mathcal{S}(M, \Omega) \) the associated operator \( \delta^* = \delta^*_{\nabla} \) satisfies:

1. \( \delta^* \) is a derivation of the symmetric product \( \odot \). That is,

\[
\delta^*(\alpha \odot \beta) = \delta^* \alpha \odot \beta + \alpha \odot \delta^* \beta \quad \text{for } \alpha \in S^k(M) \text{ and } \beta \in S^l(M).
\]

2. The Schouten bracket \([\cdot, \cdot]\) is the coboundary of \( \delta^* \) in the Lie algebra cochain complex of the algebraic bracket \((\cdot, \cdot)\) with coefficients in its adjoint representation. That is,

\[
[\alpha, \beta] = (\delta^* \alpha, \beta) + (\alpha, \delta^* \beta) - \delta^* (\alpha, \beta) \quad \text{for } \alpha \in S^k(M) \text{ and } \beta \in S^l(M).
\]  

(3.9)

Consequently, the Schouten bracket \([\cdot, \cdot]\) and the algebraic Lie bracket \((\cdot, \cdot)\) are compatible Lie brackets on \( \tilde{S}(M) \). Explicitly this means

\[
0 = \text{Cycle}_{\alpha, \beta, \gamma} \left( [[\alpha, \beta], \gamma] + [(\alpha, \beta), \gamma] \right)
= ([\alpha, \beta], \gamma) + ([\beta, \gamma], \alpha) + ([\gamma, \alpha], \beta) + [[\alpha, \beta), \gamma] + [(\beta, \gamma), \alpha] + [(\gamma, \alpha), \beta].
\]  

(3.10)
3. For $\nabla = \nabla + \Pi ij^k \in \mathcal{S}(M, \Omega)$ the associated operator $\delta^*_{\nabla}$ is cohomologous to $\delta^*_\nabla$ in the Lie algebra cochain complex of the algebraic bracket $(\cdot, \cdot)$.

**Proof.** The identity (1) follows from a straightforward computation. The identity (2) follows by comparing (3.7) with the identities

$$
\delta^*(\alpha, \beta)_{i_1 \cdots i_{k+1}} = -kl(\alpha p_{i_1 \cdots i_{k-1}} \nabla_i \beta_{i_{k+1} \cdots i_{k+1}}) - \beta p_{i_1 \cdots i_{k-1}} \nabla_i \alpha_{i_{k+1} \cdots i_{k+1}},
$$

$$
(\alpha, \delta^* \beta)_{i_1 \cdots i_{k+1}} = -k\alpha p_{i_1 \cdots i_{k-1}} \nabla_i \beta_{i_{k+1} \cdots i_{k+1}} - kl\alpha p_{i_1 \cdots i_{k-1}} \nabla_i \beta_{i_{k+1} \cdots i_{k+1}},
$$

which follow straightforwardly from the definition (3.1). Since $[\cdot, \cdot]$ is a coboundary with respect to $(\cdot, \cdot)$ it is automatically compatible with $(\cdot, \cdot)$. By claim (2), $\delta^*_{\nabla} - \delta^*_\nabla$ is closed in the Lie algebra cochain complex of the algebraic bracket $(\cdot, \cdot)$. For $\Pi \in \mathcal{T}_\nabla \mathcal{S}(M, \Omega)$, $\nabla = \nabla + \Pi, \nabla \in \mathcal{S}(M, \Omega)$, and $\gamma \in \mathcal{S}_k(M)$, there holds

$$
\delta^*_{\nabla} \gamma_{i_1 \cdots i_{k+1}} = \delta^*_\nabla \gamma_{i_1 \cdots i_{k+1}} + \frac{1}{k} (\gamma, \Pi)_{i_1 \cdots i_{k+1}}.
$$

The identity (3.11) means that $\delta^*_{\nabla} - \delta^*_\nabla$ is exact, that is, that $\delta^*_\nabla$ and $\delta^*_{\nabla}$ are cohomologous in the Lie algebra cochain complex of the algebraic bracket $(\cdot, \cdot)$.

**Lemma 3.6.** For $s, t \in \mathbb{R}^t$ let $[\cdot, \cdot]_{s,t} = s(\cdot, \cdot) + t[\cdot, \cdot]$. If $s, t, s, t$, and $t$ are nonzero, the Lie brackets $[\cdot, \cdot]_{s,t}$ and $[\cdot, \cdot]_{s,t}$ are isomorphic.

**Proof.** For $m \in \mathbb{Z}$ and $t \in \mathbb{R}^t$, define $\Theta^m_t : \mathcal{S}(M) \to \mathcal{S}(M)$ by $\Theta^m_t (\alpha)_k = t^{k-m} \alpha_k$. The maps $\Theta^m_t \in \text{End}(\mathcal{S}(M))$ are linear isomorphisms such that $\Theta^0_t$ preserves $(\cdot, \cdot)$ and $\Theta^1_t$ preserves $[\cdot, \cdot]$. For $t, s \in \mathbb{R}^t$, the linear isomorphism $\Psi_{s,t} = \Theta^2_{s-t} \circ \Theta^1_t$ satisfies $[\Psi_{s,t}(\alpha), \Psi_{s,t}(\beta)]_{1,1} = \Psi_{s,t}([\alpha, \beta]_{s,t})$, hence $\Psi_{s,t}$ is an isomorphism from $(\mathcal{S}(M), [\cdot, \cdot]_{s,t})$ to $(\mathcal{S}(M), [\cdot, \cdot]_{1,1})$ if $s, t \in \mathbb{R}^t$.

The brackets $[\cdot, \cdot]_{s,t}$ interpolate between $(\cdot, \cdot)$ and $[\cdot, \cdot]$. Lemma 3.6 means that linear combinations of $(\cdot, \cdot)$ and $[\cdot, \cdot]$ in which both brackets figure nontrivially are all isomorphic. A convenient such linear combination is the bracket

$$
[\alpha, \beta] = (\alpha, \beta) + [\alpha, \beta].
$$

With respect to $[\cdot, \cdot]$, the subspace $\mathcal{S}_2(M)$ is a subalgebra of $\mathcal{S}(M)$, and the subspaces $\mathcal{S}_k(M) \subset \mathcal{S}_2(M)$ are Lie ideals. Since $\Psi_{s,t}$ preserves the filtration of $\mathcal{S}(M)$ by the subspaces $\mathcal{S}_k(M)$, these, and other similar, statements that depend only on the filtration, are valid also for the Lie brackets $[\cdot, \cdot]_{s,t}$.

Let $\alpha_i \in \mathcal{S}(M)$ be the projection onto $\mathcal{S}(M)$ of $\alpha \in \mathcal{S}(M)$.

**Lemma 3.7.** Let $(M, \Omega)$ be a symplectic manifold. With respect to the bracket $[\cdot, \cdot]$ of (3.12) and the symmetric product $\circ$, the subspace

$$
\mathcal{H}(M) = \{ \alpha \in \mathcal{S}(M) : \alpha_1 = -\delta^* \alpha_0 \}
$$

is a Poisson subalgebra of $(\mathcal{S}(M), [\cdot, \cdot])$. For $k \geq 0$, let $\mathcal{H}_k(M) = \{ \alpha \in \mathcal{H}(M) : \alpha_i = 0 \text{ if } i < k \}$ (so $\mathcal{H}_k(M) = \mathcal{S}_k(M)$ for $k \geq 2$). For $k \geq 0$ and $l \geq 2$, $[\mathcal{H}_k(M), \mathcal{H}_l(M)] \subset \mathcal{H}_{k+l-2}(M)$. In particular, for $k \geq 2$, $(\mathcal{H}_k(M), [\cdot, \cdot], \circ)$ is a Poisson ideal in $\mathcal{H}(M, [\cdot, \cdot], \circ)$. Define $\pi : \mathcal{H}(M) \to C^\infty(M)$ by $\pi(\alpha) = -\alpha_0$. The sequence

$$
\{0\} \to (\mathcal{H}_2(M), [\cdot, \cdot], \circ) \to (\mathcal{H}(M), [\cdot, \cdot], \circ) \to (C^\infty(M), [\cdot, \cdot]) \to \{0\}
$$

of Poisson algebras is exact and $\iota : (C^\infty(M), [\cdot, \cdot]) \to (\mathcal{H}(M), [\cdot, \cdot])$ defined by $\iota(f) = -f + \delta^* f$ is an injective Lie algebra homomorphism such that $\pi \circ \iota$ is the identity. That is, the sequence (3.14) splits and $\mathcal{H}(M)$ is the semidirect product of the Lie subalgebra

$$
\mathfrak{h}(M) = \iota(C^\infty(M)) = \{ \alpha \in \mathcal{S}(M) : \alpha_1 = -\delta^* \alpha_0, \alpha_i = 0 \text{ for } i \geq 2 \}
$$

and the Poisson ideal $\mathcal{H}_2(M)$. 


**Proof.** Combining the \( k = 0 \) case of (3.7) and the \( k = 1 \) case of (3.1), shows that
\[
[f, \alpha] = -k\iota(H_f)\alpha = (\delta^* f, \alpha), \quad \text{for } f \in C^\infty(M) \quad \text{and} \quad \alpha \in \mathcal{S}^k(M). \quad (3.15)
\]
Alternatively, (3.15) is a special case of (3.9). Let \( \alpha, \beta \in \mathbf{H}(M) \). By (3.15),
\[
(\alpha_1, \beta_{j+1}) + (\alpha_{j+1}, \beta_1) + [\alpha_0, \beta_{j+1}] + [\alpha_{j+1}, \beta_0]
= - (\delta^* \alpha_0, \beta_{j+1}) - (\alpha_{j+1}, \delta^* \beta_0) + [\alpha_0, \beta_{j+1}] + [\alpha_{j+1}, \beta_0] = 0,
\]
for all \( j \geq 0 \). By (3.16),
\[
[\alpha, \beta_0] = (\alpha_1, \beta_1) + [\alpha_0, \beta_1] + [\alpha_1, \beta_0]
= (\delta^* \alpha_0, \delta^* \beta_0) - [\delta^* \alpha_0, \delta^* \beta_0] = -(\delta^* \alpha_0, \delta^* \beta_0) = -\{\alpha_0, \beta_0\}, \quad (3.17)
\]
and, by (3.16) and (3.8),
\[
[\alpha, \beta_1] = (\alpha_1, \beta_2) + (\alpha_2, \beta_1) + [\alpha_0, \beta_2] + [\alpha_1, \beta_1] + [\alpha_2, \beta_0]
= [\delta^* \alpha_0, \delta^* \beta_0] = \delta^* \{\alpha_0, \beta_0\} = -\delta^*[\alpha, \beta]_0. \quad (3.18)
\]
By (3.17) and (3.18), \( \mathbf{H}(M), [\cdot, \cdot] \rangle \) is a Lie subalgebra of \( \mathbf{S}(M), [\cdot, \cdot] \rangle \) and \( \mathbf{H}(M) \) is a Lie ideal of \( \mathbf{H}(M), [\cdot, \cdot] \rangle \). That \( \mathbf{H}(M), [\cdot, \cdot], (\cdot) \rangle \) is a Poisson subalgebra of \( \mathbf{S}(M), [\cdot, \cdot], (\cdot) \rangle \) follows from
\[
\delta^* (\alpha_0 \beta_0) = -\alpha_0 \beta_1 - \beta_0 \alpha_1 = -(\alpha \circ \beta)_1.
\]
Let \( \alpha, \beta \in \mathbf{H}(M) \) and let \( j \geq 2 \). Then,
\[
[\alpha, \beta]_j = \sum_{i=1}^{j+1} (\alpha_i, \beta_{j+2-i}) + \sum_{i=0}^{j+1} [\alpha_i, \beta_{j+1-i}] = \sum_{i=2}^{j+1} (\alpha_i, \beta_{j+2-i}) + \sum_{i=1}^{j} [\alpha_i, \beta_{j+1-i}], \quad (3.19)
\]
where the second equality follows from (3.16). Suppose \( k \geq 0, l \geq 2, \alpha \in \mathbf{H}_k(M), \) and \( \beta \in \mathbf{H}_l(M) \).
If \( j < k + l - 2 \) and \( i \geq k \) then \( j + 1 - i < j + 2 - i \leq j + 2 - k < l \); hence every term in the sums in the last line of (3.19) vanishes. This shows that \( \mathbf{H}_k(M), \mathbf{H}_l(M) \rangle \subset \mathbf{H}_{k+l-2}(M) \). It also follows from (3.19) that \( \mathbb{H}(M) \) is a Lie subalgebra of \( \mathbf{H}(M), [\cdot, \cdot] \rangle \). By (3.17),
\[
\pi([\alpha, \beta]) = -[\alpha, \beta]_0 = (\delta^* \alpha_0, \delta^* \beta_0) = \{\alpha_0, \beta_0\} = [\pi(\alpha), \pi(\beta)],
\]
so that \( \pi \) is a Lie algebra homomorphism. For \( f, g \in C^\infty(M) \), by (3.15) and (3.8),
\[
[\iota(f), \iota(g)] = [-f + \delta^* f, -g + \delta^* g] = (\delta^* f, \delta^* g) - [f, \delta^* g] + [\delta^* f, g] = -(\delta^* f, \delta^* g) + \delta^* \{f, g\} = \iota(\{f, g\}),
\]
showing that \( \mathbb{H}(M) = \iota(C^\infty(M)) \) is a Lie subalgebra of \( \mathbf{H}(M), [\cdot, \cdot] \rangle \) that is isomorphic to \( (C^\infty(M), [\cdot, \cdot]) \).

**Lemma 3.8.** Let \( (M, \Omega) \) be a symplectic manifold. For a symplectic vector field \( X \in \text{symp}(M, \Omega) \),

1) the Lie derivative \( \mathfrak{L}_X \) along \( X \) is a derivation of the algebraic bracket \( (\cdot, \cdot) \);

2) the operator \( I(X): \mathbf{S}(M) \rightarrow \mathbf{S}(M) \) defined by \( I(X) \alpha = k\iota(X)\alpha = -(X^\alpha, \alpha) \) for \( \alpha \in \mathbf{S}^k(M) \) is a derivation of the Schouten bracket \( [\cdot, \cdot] \cdot \).

\[\square\]
Proof. Let $\alpha = X^\gamma$, $\beta \in S^k(M)$, and $\gamma \in S^l(M)$. By (3.10) and the definition of $I(X)$,

$$
\mathcal{L}_X(\beta, \gamma) - (\mathcal{L}_X(\beta), \gamma) - (\beta, \mathcal{L}_X(\gamma)) = [\alpha, \beta, \gamma] - ([\alpha, \beta], \gamma) - (\beta, [\alpha, \gamma])
$$

$$
= -(\alpha, [\beta, \gamma]) + ([\alpha, \beta], \gamma) + \beta, ([\alpha, \gamma])
$$

$$
= I(X)[\beta, \gamma] - [I(X)\beta, \gamma] - [\beta, I(X)\gamma].
$$

By (3.20), (1) and (2) are equivalent, so it suffices to prove either. As the claims are local in nature, it can be assumed $\alpha$ is exact. Taking $\alpha = -\delta^* f$, for $f \in C^\infty(M)$, in (3.20) yields, by (3.15),

$$
(\delta^* f, [\beta, \gamma]) - (\delta^* f, [\beta, \gamma]) - (\beta, (\delta^* f, \gamma)) = [f, [\beta, \gamma]] - ([f, \beta], \gamma) - ([\beta, f], \gamma) = 0.
$$

Let $\pi_k : \hat{S}(M) \to \hat{S}(M)$ be the projection with image $\hat{S}_k(M)$ defined by $\pi_k \alpha = \alpha - \sum_{i=0}^{k-1} \alpha_i$. Define truncated brackets on $\hat{S}(M)$ by $(\alpha, \beta)_{(k)} = (\pi_k \alpha, \pi_k \beta)$ and $[\alpha, \beta]_{(k)} = [\pi_k \alpha, \pi_k \beta]$. If $k \geq 2$ then $2k - 2 \geq k$ so $(\alpha, \beta)_{(k)} \in \hat{S}_k(M)$ and so $\pi_k(\alpha, \beta)_{(k)} = (\alpha, \beta)_{(k)}$; this identity suffices to show that $(\cdot, \cdot)_{(k)}$ satisfies the Jacobi identity, so is a Lie bracket on $\hat{S}(M)$. The same argument shows that if $k \geq 1$ then $\hat{S}_{k+1}(M)$ is an ideal with respect to each of the brackets $[\cdot, \cdot]_{(k)}$ and $(\cdot, \cdot)_{(k+1)}$.

Lemma 3.9. Let $(M, \Omega)$ be a symplectic manifold. For $k \geq 1$, the Lie brackets $[\cdot, \cdot]_{(k)}$ and $(\cdot, \cdot)_{(k+1)}$ on $\hat{S}_{k+1}(M)$ are compatible.

Proof. For $\alpha, \beta \in \hat{S}_{k+1}(M)$, by the definitions of $[\cdot, \cdot]_{(k)}$ and $(\cdot, \cdot)_{(k+1)}$, and (3.9),

$$
[\alpha, \beta]_{(k)} = [\pi_k \alpha, \pi_k \beta] = (\delta^* \pi_k \alpha, \pi_k \beta) + (\pi_k \alpha, \delta^* \pi_k \beta) - \delta^* (\pi_k \alpha, \pi_k \beta)
$$

$$
= (\pi_{k+1} \delta^* \alpha, \pi_{k+1} \beta) + (\pi_{k+1} \alpha, \pi_{k+1} \delta^* \beta) - \delta^* (\pi_{k+1} \alpha, \pi_{k+1} \beta)
$$

$$
= (\delta^* \alpha, \beta)_{(k+1)} + \alpha, (\delta^* \beta)_{(k+1)} - \delta^* (\alpha, \beta)_{(k+1)},
$$

where the penultimate equality follows from $\delta^* \pi_k = \pi_{k+1} \delta^*$ and $\pi_k \pi_{k+1} = \pi_{k+1}$ (since $\alpha \in \hat{S}_{k+1}(M)$, the latter implies $\pi_k \alpha = \pi_{k+1} \alpha$). This shows that $[\cdot, \cdot]_{(k)}$ is a cocycle with respect to $(\cdot, \cdot)_{(k+1)}$.

That an arbitrary vector field need not preserve $\Omega$ suggests working with the subspace $\tilde{Z}(M)$ of $\hat{S}(M)$, the degree one part of which comprises closed one-forms on $M$ (the symplectic dual $\text{symp}(M, \Omega)^\gamma$ of $\text{symp}(M, \Omega)$). (A related problem is that the $(\cdot, \cdot)$-bracket of a closed one-form with an element of $S^2(M)$ need not be a closed one-form.) Define

$$
\tilde{B}(M) = \{ \alpha \in \hat{S}(M) : \alpha_1 \text{ is exact} \} \subset \tilde{Z}(M) = \{ \alpha \in \hat{S}(M) : d\alpha = 0 \}.
$$

If $\alpha, \beta \in \tilde{Z}(M)$ then $[\alpha, \beta]_{(1)} = [\pi_1 \alpha, \pi_1 \beta] = [\alpha_1, \beta_1]$ is closed, since the Schouten bracket of closed one-forms is closed. Hence $\tilde{Z}(M)$ is a subalgebra of $\hat{S}(M)$ with respect to the bracket $[\cdot, \cdot]_{(1)}$. If $\alpha, \beta \in \tilde{Z}(M)$ then $(\alpha, \beta)_{(2)} \in \tilde{S}_2(M) = \tilde{Z}(M) \cap \tilde{S}_2(M)$, so $\tilde{Z}(M)$ is a subalgebra of $\hat{S}(M)$ with respect to the bracket $(\cdot, \cdot)_{(2)}$.

Lemma 3.10. Let $(M, \Omega)$ be a symplectic manifold. The Lie brackets $[\cdot, \cdot]_{(1)}$ and $(\cdot, \cdot)_{(2)}$ on $\tilde{Z}(M)$ are compatible and $\tilde{B}(M)$ is an ideal in $\tilde{Z}(M)$ with respect to each of $[\cdot, \cdot]_{(1)}$ and $(\cdot, \cdot)_{(2)}$.

Proof. Let $\alpha, \beta, \gamma \in \tilde{Z}(M)$. Then, using $\pi_1 \alpha = \alpha_1 + \pi_2 \alpha$,

$$
\text{Cycle}(([\alpha, \beta]_{(1)}, \gamma)_{(2)} + ([\alpha, (\beta)_{(2)}, \gamma]_{(1)})
$$
the last equality because, by Lemma 3.8, the Lie derivative along a symplectic vector field is a derivation of the algebraic bracket (·, ·). This shows \([·, ·]_{(1)}\) and \((·, ·)_{(2)}\) are compatible. If \(α ∈ \hat{B}(M)\) and \(β ∈ \hat{Z}(M)\), then \((α, β)_{(2)} ∈ \hat{S}_2(M) ⊂ \hat{B}(M)\), and there is \(f ∈ C^∞(M)\) such that \(α_1 = df\), so that
\[
[α_1, β_1]_1 = [π_1α_1, π_1β_1]_1 = [α_1, β_1] = [β_1, df] = \mathcal{L}_{β_1} df = d\mathcal{L}_{β_1} f.
\]
Hence \([α, β] ∈ \hat{B}(M)\). □

By Lemma 3.10 the bracket
\[
[·, ·]_τ = (·, ·)_{(2)} + [·, ·]_{(1)}
\]
(the subscript \(τ\) is meant to suggest truncated) is a Lie bracket on \(\hat{Z}(M)\) and \(\hat{B}(M)\) is an ideal of \((\hat{Z}(M), [·, ·]_τ)\). Because \([α, β]_τ ∈ \hat{S}_1(M)\) for all \(α, β ∈ \hat{Z}(M)\), \(\hat{Z}_1(M) = \hat{Z}(M) ∩ \hat{S}_1(M)\) and \(\hat{B}_1(M) = \hat{B}(M) ∩ \hat{S}_1(M)\) are ideals in \((\hat{Z}(M), [·, ·]_τ)\). Lemma 3.11 shows that the Lie algebra \((\hat{H}(M), [·, ·])\) is a central extension of \((\hat{B}_1(M), [·, ·]_τ)\).

**Lemma 3.11.** Let \((M, Ω)\) be a symplectic manifold. The sequence
\[
\{0\} → \mathbb{R} → \hat{H}(M), [·, ·] → \hat{B}_1(M), [·, ·]_τ → \{0\}
\]
(3.23)
of Lie algebras is exact, where \(i(c)\) is the constant function on \(M\) equal to \(-c\).

**Proof.** For \(α ∈ \hat{H}(M)\), \(π_1α = 0\) if and only if \(α_i = 0\) for \(i ≥ 2\) and \(α_0\) is constant. As \(π_1\) maps \(\hat{H}(M)\) surjectively onto \(\hat{B}_1(M)\), the content of the claim is that \(π_1\) is a Lie algebra homomorphism. Let \(α, β ∈ \hat{H}(M)\). The identity (3.19) can be restated as
\[
π_2[α, β] = π_2[π_1α, π_1β] + π_2[π_1α_1, π_1β_1] = π_2[α, β]_τ.
\]
The identity (3.18) shows \([α, β]_1 = δ^*\{α_0, β_0\}\). On the other hand,
\[
[α, β]_1 = [π_2α_1, π_2β_1] + [π_1α_1, π_1β_1] = [α_1, β_1]
\]
\[
= [δ^*α_0, δ^*β_0] = δ^*\{α_0, β_0\} = [α, β]_1.
\]
Hence \(π_1[α, β] = π_1[α, β]_τ = [α, β]_τ = [π_1α, π_1β]_τ\). □

The subspace \(\hat{S}_2(M)\) is a Lie ideal of each of \((\hat{H}(M), [·, ·])\) and \((\hat{B}(M), [·, ·]_τ)\), and the restrictions to \(\hat{S}_2(M)\) of the brackets \([·, ·]\) and \([·, ·]_τ\) coincide.
4 Functionals on the symplectic affine space \((S(M, \Omega), \Omega)\) of symplectic connections

This section reviews the definitions and characterizations of the preferred and critical symplectic connections.

4.1. The Lie derivative of the affine connection \(\nabla\) along \(X\),

\[
(\mathcal{L}_X \nabla)_{ij}^k = \nabla_i \nabla_j X^k + X^p R_{pij}^k
\]

is defined to be the derivative at \(t = 0\) of the pullback \(\phi_t^* \nabla\) along the flow \(\phi_t\) of \(X\). Given \(\nabla \in S(M, \Omega)\), define \(\mathcal{L} = \mathcal{L}_\nabla : \Gamma(T^*M) \to T\nabla S(M, \Omega)\) by \(\mathcal{L}(X^\gamma)_{ijk} = (\mathcal{L}_X \nabla)_{ij}^k\) for \(X \in \Gamma(TM)\). Since \(2(\mathcal{L}_X \nabla)_{ijk} = \nabla_i \nabla_j X^\gamma + \frac{2}{3} \nabla_k (dX^\gamma)_{ijk}\),

\[
(\mathcal{L}_X \nabla)_{ijk} = (\mathcal{L}_X \nabla)_{(ijk)} + \frac{2}{3} \nabla_k (dX^\gamma)_{ijk} = \mathcal{L}(X^\gamma)_{ijk} + \frac{2}{3} \nabla_k (dX^\gamma)_{ijk}.
\]

If \(X \in \text{symp}(M, \Omega)\), then \((\mathcal{L}_X \nabla)_{ijk}\) is completely symmetric, so, in this case, \(\mathcal{L}(X^\gamma)_{ijk} = (\mathcal{L}_X \nabla)_{ijk}\). Define a linear operator \(\mathcal{S} : S^3(M) \to \Gamma(T^*M)\) by \(\mathcal{S}(\beta)_{ij} = \beta^{pde} R_{i j k}\). The formal adjoint \(\mathcal{S}^* : \Gamma(T^*M) \to S^3(M)\) is \(\mathcal{S}^* (\alpha)_{ijk} = -\alpha^p R_p(i j k)\). For \(\alpha \in \Gamma(T^*M)\), \(\mathcal{L}(\alpha)\) can be rewritten as

\[
\mathcal{L}(\alpha)_{ijk} = \nabla_i \nabla_j \alpha_k + \alpha^p R_p(i j k) = (\delta^{*2} \alpha)_{ijk} - S^*(\alpha)_{ijk}.
\]

The induced action of \(\text{ham}(M, \Omega)\) on \(S(M, \Omega)\) is given by \(\mathcal{H} : C^\infty(M) \to T\nabla S(M, \Omega)\) defined by

\[
\mathcal{H}(f) = \mathcal{L}_H f = \mathcal{L}(-df) = \mathcal{L}(\delta^* f) = (\delta^{*3} f) - S^*(\delta^* f).
\]

Taking formal adjoints shows \(\mathcal{L}^* = -\delta^* - S\) and \(\mathcal{H}^* = (\mathcal{L}\delta^*)^* = \delta\mathcal{L}^* = -\delta^3 - \delta S\). Tracing (4.1) and using the Ricci identity twice yields that, for \(\nabla \in S(M, \Omega)\),

\[
\delta \mathcal{L}(X^\gamma)_{ij} = \nabla_p (\mathcal{L}_X \nabla)_{ij}^p = \nabla_p \nabla_i \nabla_j X^p + X^q \nabla_p R_{qij}^p + R_{qij}^p \nabla_p X^q
\]

\[
= \nabla_i \nabla_p \nabla_j X^p + R_{iq} \nabla_j X^q + 2X^q \nabla_q R_{ij} + \nabla_j \nabla_p \nabla_i X^q + R_{iq} \nabla_j X^q + X^q \nabla_q R_{ij} + 2X^q \nabla_j R_{iq}
\]

\[
= (\mathcal{L}_X \text{Ric})_{ij}.
\]

Taking \(X = H_f\) in (4.2) yields \(\delta \mathcal{H}(f) = \mathcal{L}_H f\).

4.2. The Cahen–Gutt moment map \(\mathcal{K} : S(M, \Omega) \to C^\infty(M)\) is defined by

\[
\mathcal{K}(\nabla) = -\delta^2 \text{Ric} - \frac{1}{2} R_{ij} R_{ij} + \frac{1}{4} R_{ijkl} R_{ijkl} = \nabla^i \nabla_j R_{ij} - \frac{1}{2} R_{ij} R_{ij} + \frac{1}{4} R_{ijkl} R_{ijkl}.
\]

The map \(\mathcal{K}\) was defined by M. Cahen and S. Gutt in [7] (see also [2] or [13]), where it was shown that \(\mathcal{K}\) is a moment map for the action of \(\text{Ham}(M, \Omega)\) on \(S(M, \Omega)\). The map \(\mathcal{K}(\nabla)\) has the form required by Tamarkin’s theorem; its part quadratic in the curvature is a constant multiple of the contraction of the first Pontryagin form \(p_1(\nabla)\) of \(\nabla\) with \(\Omega_2\): \(\mathcal{K}(\nabla) = -\delta^2 \text{Ric} - \frac{1}{2} p_1(\nabla) p_1\). As a consequence, when \(M\) is compact the integral of \(\mathcal{K}(\nabla)\) depends only on the cohomology class \([\Omega_n]\) of \(\Omega_n\) and the first Pontryagin class \([p_1]\) of \(M\),

\[
\int_M \mathcal{K}(\nabla) \Omega_n = -4\pi^2 \int_M p_1 \wedge \Omega_{n-2} = -4\pi^2 (|p_1| \cup [\Omega_{n-2}], [M]).
\]

(See [10] for details.) Since integration against \(f\Omega_n\) defines a linear functional on \(\text{ham}(M, \Omega)\) for any \(f \in C^\infty(M)\), whatever is the precise meaning of the dual space \(\text{ham}(M, \Omega)^*\), this
space contains $\text{ham}(M, \Omega)$ as a subspace. Hence $\mathcal{K}$ can be regarded as a map $\mathcal{K}: \mathcal{S}(M, \Omega) \to \text{ham}(M, \Omega)^*$.

The first variation $\delta_{\Pi} F(\nabla)$ of a functional $F$ on $\mathcal{S}(M, \Omega)$ at $\nabla \in \mathcal{S}(M, \Omega)$ in the direction of $\Pi \in T_{\nabla} \mathcal{S}(M, \Omega)$ is defined by $\delta_{\Pi} F(\nabla) = \frac{d}{dt}|_{t=0} F(\nabla + t\Pi)$. The essential content of Theorem 4.1 is that $\delta_{\Pi} \mathcal{K}_\nabla = \mathcal{H}_\nabla(\Pi)$.

**Theorem 4.1** (M. Cahen and S. Gutt, [7]). Let $(M, \Omega)$ be a symplectic manifold. The map $\mathcal{K}: \mathcal{S}(M, \Omega) \to \text{ham}(M, \Omega)^*$ is a moment map for the action of $\text{Ham}(M, \Omega)$ on $\mathcal{S}(M, \Omega)$, equivariant with respect to the actions of $\text{Symp}(M, \Omega)$.

**Indication of proof.** For a complete proof see [10]. Let $\tilde{R}_{ijk}^l$ be the curvature of $\nabla = \nabla + t\Pi_{ij}^k$ and label with a * the tensors derived from it; for example $\tilde{R}_{ij}$ is the Ricci curvature of $\nabla$. Let $B(\Pi)_{ij} = \Pi_{ip}^q \Pi_{jq}^p$ and $(\Pi^*)_{ii} = 3\delta B(\Pi)_{ii} - \Pi^{abc} \nabla_i \Pi_{abc}$. Then

$$\begin{align*}
\tilde{R}_{ijkl} &= R_{ijkl} + 2t \nabla_i \Pi_{jkl} + 2t^2 \Pi_{ip}^r \Pi_{jkl} \delta^p_r = R_{ijkl} + td \nabla_i \Pi_{ijkl} + t^2 (\Pi \wedge \Pi)_{ijkl}, \\
\tilde{R}_{ij} &= R_{ij} + t \nabla_i \Pi_{ij} + t^2 \Pi_{ip}^q \Pi_{jq}^p = R_{ij} + td \nabla_i \Pi_{ij} - t^2 B(\Pi)_{ij}, \\
\delta_{\nabla} \tilde{R}_{ij} &= \delta_{\nabla} R_{ij} + t(\delta^* \Pi + \frac{1}{3} (\text{Ricci}(\Pi))) + t^2 (\delta^* B(\Pi) + \frac{1}{3} (\delta \Pi, \Pi)) - \frac{1}{3} \delta^3 (B(\Pi), \Pi), \\
\delta_{\nabla} + t \Pi \nabla_i \text{Ricci}(\nabla + \Pi)_{ii} &= \delta_{\nabla} \text{Ricci}_{ii} - t (\mathcal{L}^* (\Pi) + \mathcal{T}^* (\Pi)) - t^2 (\delta \Pi, B(\Pi)) + \Pi_{ij} \Pi_{pq} + O(t^3).
\end{align*}$$

where $\mathcal{T}(\beta)_{ii} = \beta_{abc} (R_{iabc} - \Omega_{(a R_{bc})}) = \beta_{ij}^p R_{pq} + \delta(\beta)_{ii}$. Combining (4.5)–(4.8) yields

$$\mathcal{K}(\nabla + \Pi) = \mathcal{K}(\nabla) + t \mathcal{H}^* (\Pi) + \frac{1}{2} t^2 \delta (\Pi^* \Pi) + O(t^3).$$

For $\phi \in \text{Symp}(M, \Omega)$, $\mathcal{K}(\phi^* \nabla) = \mathcal{K}(\nabla) \circ \phi$, so $\langle f \circ \phi, \mathcal{K}(\phi^* \nabla) \rangle = \langle f, \mathcal{K}(\nabla) \rangle$. This shows $\mathcal{K}$ is equivariant with respect to the action of $\text{Symp}(M, \Omega)$ on $\mathcal{S}(M, \Omega)$ and its action on $\text{ham}(M, \Omega)^*$ induced by its action on $C^\infty(M)$. For $\nabla \in \mathcal{S}(M, \Omega)$, $\Pi \in T_{\nabla} \mathcal{S}(M, \Omega)$, $\chi \in \text{symp}(M, \Omega)$, and $f \in C^\infty(M)$ it follows from (4.9) that $\delta_{\Pi} \mathcal{K}_{\nabla} = \mathcal{H}_{\nabla}(\Pi)$. Hence $\mathcal{K}_{\nabla}(\mathcal{H}(f), \Pi) = \langle f, \mathcal{H}^* (\Pi) \rangle = \delta_{\Pi} \langle f, \mathcal{K}(\nabla) \rangle$, showing that $\mathcal{K}$ is a moment map.

**Lemma 4.2.** For any $f \in C^4(\mathbb{R})$, the Hamiltonian vector field on $(\mathcal{S}(M, \Omega), \Omega)$ generated by $\mathcal{E}_f(\nabla) = \int_M f(\mathcal{K}(\nabla)) \Omega_n$ is $\mathcal{H}_{\nabla} = -\mathcal{H}(f'(\mathcal{K}(\nabla)))$.

**Proof.** Calculating the first variation $\delta_{\Pi} \mathcal{E}_f(\nabla)$ along $\Pi \in T_{\nabla} \mathcal{S}(M, \Omega)$ using (4.9) yields

$$\delta_{\Pi} \mathcal{E}_f = \langle f'(\mathcal{K}(\nabla)), \mathcal{H}^* (\Pi) \rangle = \Omega_{\nabla} \langle \mathcal{H}(f'(\mathcal{K}(\nabla))), \Pi \rangle,$$

from which the claim follows.

A symplectic connection $\nabla \in \mathcal{S}(M, \Omega)$ is critical if it is a critical point, for arbitrary compactly supported variations $\alpha \in \mathcal{S}^3(M)$, of $\mathcal{E}: \mathcal{S}(M, \Omega) \to C^\infty(M)$ defined by $\mathcal{E}(\nabla) = \int_M \mathcal{K}(\nabla)^2 \Omega_n$. Motivation for studying critical connections is given in [10], where the author proposed this notion.

**Corollary 4.3** ([10]). A symplectic connection $\nabla \in \mathcal{S}(M, \Omega)$ is critical if and only if $\mathcal{H}(\mathcal{K}(\nabla)) = 0$.

**4.3.** Given $\nabla \in \mathcal{S}(M, \Omega)$, there are written $\mathcal{R}_{ij}^k$ for the Ricci endomorphism $R_{ij}^k$ and $\mathcal{R}_{ij}^{ok}$ for its $k$th power as a fiberwise endomorphism. Since $(\mathcal{R}_{ij}^{ok})_{ij} = (-1)^{k+1} (\mathcal{R}_{ij}^{ok})_{ji}$, $(\mathcal{R}_{ij}^{ok})_{ij}^k$ is symmetric if $k$ is odd, and skew-symmetric if $k$ is even. In particular, $\text{tr} \mathcal{R}_{ij}^{ok} = 0$ and

$$\text{tr} \mathcal{R}_{ij}^{ok} = -R_{ij}^k (\mathcal{R}_{ij}^{ok})_{ij} = (-1)^k R_{1ij2} R_{2ij3} \cdots R_{2k-2i2k-1} R_{2k-1i2k} R_{2k1i}.$$

**Remark 4.4.** The functionals $\text{tr} \mathcal{R}$ play a key role in the averaging procedure used in [9, Section 2], where they are called *cycles of length s*.
For $1 \leq k$ and $\text{Ric}^{\circ k}_{ij} = (\text{Rc}^{\circ k})_{ij}$, define $\mathcal{R}(k): \mathbb{S}(M, \Omega) \to \mathbb{R}$ by

$$
\mathcal{R}(k)(\nabla) = -\frac{1}{2k} \int_M \text{tr} \left( \text{Rc}^{\circ 2k-1} \right) \Omega_n
= \frac{1}{2k} \int_M R^{pq} (\text{Rc}^{\circ 2k-1})_{pq} \Omega_n
= \frac{1}{2k} \int_M R^{pq} (\text{Rc}^{\circ 2k-1})_{pq} \Omega_n
= \frac{1}{2k} \int_M \nabla \left( \text{Rc}^{\circ 2k-1}, \text{Ric} \right).
$$

(4.11)

**Lemma 4.5.** The Hamiltonian vector field $\mathcal{H}_{\mathcal{R}(k)}$ on $(\mathbb{S}(M, \Omega), \Omega)$ generated by $\mathcal{R}(k)$ is

$$
\mathcal{H}_{\mathcal{R}(k)}(\nabla)_{abc} = -\delta^* \text{Ric}^{\circ 2k-1}_{abc} = \nabla(a (\text{Rc}^{\circ 2k-1})_{bc}).
$$

**Proof.** For $\Pi \in T_{\nabla} \mathbb{S}(M, \Omega)$, by (4.6) and (4.11),

$$
\Omega_{\nabla}(\Pi, \mathcal{H}_{\mathcal{R}(k)}) = \delta \Pi \mathcal{R}(k)(\nabla) = -\frac{1}{2k} \delta \Pi \int_M \text{tr} \left( \text{Rc}^{\circ 2k} \right) \Omega_n
= \frac{1}{2k} \delta \Pi \int_M R^{pq} (\text{Rc}^{\circ 2k-1})_{pq} \Omega_n
= \langle \delta \Pi, \text{Ric}^{\circ 2k-1} \rangle
= -\Omega_{\nabla}(\Pi, \delta^* \text{Ric}^{\circ 2k-1}).
$$

(4.12)

**Remark 4.6.** Since, by definition, $\nabla \in \mathbb{S}(M, \Omega)$ is preferred if it is critical for $\mathcal{R}(1)$, that is a zero of $\mathcal{H}_{\mathcal{R}(1)}$, Lemma 4.5 suggests regarding the zeros of $\mathcal{H}_{\mathcal{R}(k)}$, for $k > 1$, as generalizations of preferred symplectic connections.

**Remark 4.7.** For any $\Psi \in \text{Symp}(M, \Omega)$, it follows from $\mathcal{K}(\Psi^* \nabla) = \mathcal{K}(\nabla) \circ \Psi$ and the similar equivariance of the Ricci tensor, that $\mathcal{E}_\phi(\Psi^* \nabla) = \mathcal{E}_\phi(\nabla)$ and $\mathcal{R}(k)(\Psi^* \nabla) = \mathcal{R}(k)(\nabla)$. Consequently, $\mathcal{E}_\phi$ and $\mathcal{R}(k)$ are constant on $\text{Symp}(M, \Omega)$ orbits in $\mathbb{S}(M, \Omega)$.

The Poisson bracket $\{\mathcal{F}, \mathcal{G}\}$ of functionals $\mathcal{F}$ and $\mathcal{G}$ on $(\mathbb{S}(M, \Omega), \Omega)$ is defined to be the symplectic pairing

$$
\{\mathcal{F}, \mathcal{G}\} = \Omega(\mathcal{H}_\mathcal{F}, \mathcal{H}_\mathcal{G})
$$

(4.13)

of the Hamiltonian vector fields $\mathcal{H}_\mathcal{F}$ and $\mathcal{H}_\mathcal{G}$ generated by $\mathcal{F}$ and $\mathcal{G}$.

**Lemma 4.8.** If $(M, \Omega)$ is compact, the functionals $\mathcal{R}(k)$ and $\mathcal{E}_\phi$ on $(\mathbb{S}(M, \Omega), \Omega)$ Poisson commute for any $k \geq 1$ and any $\phi \in C^1(\mathbb{R})$.

**Proof.** By definition of the Poisson bracket $\{\cdot, \cdot\}$ on $(\mathbb{S}(M, \Omega), \Omega)$, (4.2), and Lemma 4.2,

$$
\{\mathcal{E}_\phi, \mathcal{R}(k)\} = \Omega(\mathcal{H}_{\mathcal{E}_\phi}, \mathcal{H}_{\mathcal{R}(k)})
= \Omega(\mathcal{H}(\phi(\mathcal{K})), \delta^* \text{Ric}^{\circ 2k-1})
= -\langle \delta \mathcal{H}(\phi'(\mathcal{K})), \text{Ric}^{\circ 2k-1} \rangle
= -\langle \mathcal{L}_{\mathcal{H}_\phi(\mathcal{K})} \text{Ric}, \text{Ric}^{\circ 2k-1} \rangle
= \int_M (\mathcal{L}_{\mathcal{H}_\phi(\mathcal{K})} \text{Ric}) \cdot (\text{Rc}^{\circ 2k-1}) \Omega_n
= \frac{1}{2k} \int_M \mathcal{L}_{\mathcal{H}_\phi(\mathcal{K})} (\text{Rc}^{\circ 2k}) \Omega_n
= 0,
$$

the last equality because the integral of a divergence vanishes.

(4.14)

**Remark 4.9.** It would be interesting to know what can be said about existence of solutions to the Hamiltonian flows

$$
\frac{d}{dt} \nabla(t) = c \mathcal{H}_{\mathcal{R}(k)}(\nabla(t)).
$$

(4.14)

Here $\nabla(t)$ is a path in $\mathbb{S}(M, \Omega)$ and $c \in \mathbb{R}$. By Lemma 4.8, $\mathcal{E}$ is constant along the flows (4.14).
5 A symplectic affine action of \((\mathcal{H}(M), [\cdot, \cdot])\) on \((\mathcal{S}(M, \Omega), \Omega)\)

5.1. Recall from Remark 3.3 that \(\mathcal{G}(\mathcal{F})\) is the group of gauge transformations of the linear symplectic frame bundle \(\mathcal{F} \to M\). The action of \(\mathcal{G}(\mathcal{F})\), by pullback, on principal connections induce the right action, \(\mathcal{T}(M, \Omega) \times \mathcal{G}(\mathcal{F}) \to \mathcal{T}(M, \Omega)\), on the associated covariant derivatives given by

\[
g \cdot \nabla = \nabla + (g^{-1})^{k}_{p}\nabla_{i}g_{j}^{k} = \nabla - g^{bp}\nabla_{i}g_{jp}, \tag{5.1}
\]

where \(g_{iq}g_{j}^{q} = g_{i}^{p}g_{j}^{q}\Omega_{pq} = \Omega_{ij}\) and \((g^{-1})_{ij} = -g_{ji}\). The covariant derivatives associated with principal connections on \(\mathcal{F}\) preserve the symplectic structure, but need not be torsion-free, and the action of \(\mathcal{G}(\mathcal{F})\) does not preserve the torsion, for, by (5.1), if \(\nabla\) is torsion-free, then \(\nabla^{g}\) is torsion-free if and only if \(\nabla_{ij}^{g} = 0\). This need not be true in general. The action (5.1) is affine (as in Section 2.4). Its linear part \(L(g): T_{\nabla} \mathcal{T}(M, \Omega) \to T_{\nabla} \mathcal{T}(M, \Omega)\) is \((L(g))_{ijk} = (g^{-1})^{k}_{p}g_{iq}g_{j}^{p} = g_{j}^{p}g_{k}g_{ipq}\), and there follows \(\Omega(L(g)\alpha, L(g)\beta) = \Omega(\alpha, \beta)\), showing that the action of \(\mathcal{G}(\mathcal{F})\) on \(\mathcal{T}(M, \Omega)\) is symplectic affine. However, there is no obvious way to modify the action of \(\mathcal{G}(\mathcal{F})\) on \(\mathcal{T}(M, \Omega)\) to obtain an action on \(\mathcal{S}(M, \Omega)\).

5.2. The curvature is a moment map for the gauge group action on the space of principal connections on a principal bundle over a symplectic manifold. This goes back to [1], where this construction is used in the special case of a surface. In the present context this can be realized concretely as follows. Given \(\alpha \in S^{2}(M)\), regarded as an element of the Lie algebra \(\mathcal{L}(\mathcal{F})\) of infinitesimal gauge transformations (see Remark 3.3), regard the curvature of \(\nabla \in \mathcal{T}(M, \Omega)\) as a \(\mathcal{L}(\mathcal{F})\)-valued two-form \(C\) and write \(\text{tr}(C\alpha) = R_{ij}^{p}g_{iq}g_{jp}\) for the two-form resulting from tracing the composition of the curvature with the endomorphism \(\alpha_{ij}\). However the dual \(\mathcal{L}(\mathcal{F})^{*}\) is understood, it includes \(\text{Ad}(\mathcal{F})\)-valued 2n-forms.

**Theorem 5.1.** On a compact symplectic manifold \((M, \Omega)\), the map \(\Gamma: \mathcal{T}(M, \Omega) \to \mathcal{L}(\mathcal{F})^{*}\) defined by \(\Gamma(\nabla) = C \wedge \Omega_{n-1}\) is a moment map for the action of the gauge group \(\mathcal{G}(\mathcal{F})\) on \(\mathcal{T}(M, \Omega)\).

**Proof.** By (2.1), \(\Gamma(\nabla) = \frac{1}{2}R_{ij}^{p}g_{ipq}\Omega_{n}\), so

\[
\langle \Gamma(\nabla), \alpha \rangle = \int_{M} \text{tr}(C\alpha) \wedge \Omega_{n-1} = -\frac{1}{2} \int_{M} \alpha^{ij}R_{ij}^{p}g_{ipq}\Omega_{n}, \tag{5.2}
\]

If \(\nabla\) is a connection preserving a volume form \(\mu\) and having torsion \(\tau_{ij}^{k}\), then the divergence \(\nabla_{p}X^{p}\) defined by the connection and the divergence \((\Sigma X^{p})/\mu\) defined by the volume form are related by \((\Sigma X^{p})/\mu = \nabla_{p}X^{p} + X^{p}\tau_{pq}^{q}\). For \(\nabla \in \mathcal{T}(M, \Omega)\) there results

\[
\int_{M} (\nabla_{p}X^{p} + X^{p}\tau_{pq}^{q}) \Omega_{n} = \int_{M} \Sigma X^{p}\Omega_{n} = 0. \tag{5.3}
\]

The first variation of the curvature of \(\nabla \in \mathcal{T}(M, \Omega)\) along \(\Pi \in T_{\nabla} \mathcal{T}(M, \Omega)\) is \((\delta_{\Pi}R)_{ij}^{k} = d\nabla\Pi_{ijk} + \tau_{ij}^{k}R_{pkj}\). Note that \(0 = d\Omega_{ij} - 3\nabla_{[i}\Omega_{jk]} = \tau_{ij}^{k}\). Tracing this yields \(2\tau_{ip}^{p} = -\tau_{ip}^{p}\). Viewing \(\alpha_{ij} \in S^{2}(M)\) as an element of \(\mathcal{L}(\mathcal{F})\) and differentiating the action of \(\exp(\alpha)\) on \(\mathcal{T}(M, \Omega)\) yields the vector field \(X^{o}(\nabla) = \frac{d}{dt}_{t=0}(\exp(-t\alpha)\cdot \nabla) = -\nabla_{i}\alpha_{jk}\). Applying the preceding observations to (5.2) and integrating by parts using (5.3) yields

\[
\delta_{\Pi} \langle \Gamma(\nabla), \alpha \rangle = -\frac{1}{2} \delta_{\Pi} \int_{M} \alpha^{ij}R_{ij}^{p}g_{ipq}\Omega_{n} = \int_{M} (\alpha^{ij}\nabla^{p}\Pi_{pij} - \frac{1}{2} \tau_{q}^{ap}\Pi_{pij}\alpha_{ij}) \Omega_{n}
\]

\[
= \int_{M} (\Pi^{ijk}\nabla_{i}\alpha_{jk} - \nabla_{p}(\alpha^{ij}\Pi_{pij}) - \alpha^{ij}\Pi_{ij}^{pq}q) \Omega_{n}
\]

\[
= \int_{M} \Pi^{ijk}\nabla_{i}\alpha_{jk} \Omega_{n} = -\Omega(\nabla^{o}, \Pi).
\]
This shows that $X^\alpha$ is the Hamiltonian vector field on $\mathbb{T}(M, \Omega)$ generated by $\langle \Gamma(\nabla), \alpha \rangle$. Since $\Gamma$ is by construction $\mathcal{G}(\mathcal{F})$-equivariant, this shows $\Gamma$ is a moment map. ■

**Lemma 5.2.** For $\alpha \in S^2(M)$, define $R_\alpha : S(M, \Omega) \to \mathbb{R}$ by $R_\alpha(\nabla) = \langle \alpha, \text{Ric}(\nabla) \rangle$. For a compactly supported $\alpha \in S^2(M)$ the Hamiltonian vector field generated on $(S(M, \Omega), \Omega)$ by $R_\alpha$ is $H_{R_\alpha} = -\delta^* \alpha$.

**Proof.** By (4.6), $\Omega_{\nabla}(\Pi, H_{R_\alpha}) = \delta_\Pi R_\alpha(\nabla) = \Omega_{\nabla}(\alpha, \delta_\Pi \Pi) = -\Omega_{\nabla}(\Pi, \delta_\Pi^* \alpha)$. ■

If $\nabla \in S(M, \Omega)$, then $\Gamma(\nabla) = (1/2)R_p p_{ij} \Omega_n = R_{ij} \Omega_n$, because $R_p p_{ij} = 2R_{ij}$. (Note that the identity $R_p p_{ij} = 2R_{ij}$ fails for $\nabla \in \mathbb{T}(M, \Omega)$ having nonvanishing torsion.) Hence, by (5.2), $\langle \Gamma(\nabla), \alpha \rangle = -\langle \alpha, \text{Ric}(\nabla) \rangle$ for $\nabla \in S(M, \Omega)$. That is the restriction of $-\langle \Gamma(\nabla), \alpha \rangle$ to $S(M, \Omega)$ equals the functional $R_\alpha$ defined in Lemma 5.2. For $\Pi \in \Gamma(T^* M \otimes S^2(T^* M))$ let $Q(\Pi)_{ijk} = \Pi_{(ijk)}$ and $Q^2(\Pi)_{ijk} = (\text{Id} - Q)(\Pi)_{ijk} = \frac{2}{3}(\Pi_{(ijk)} - \Pi_{(jki)})$ be the projections onto $S^3(M)$ and its complement. The Hamiltonian vector field $H_{R_\alpha}$ on $S(M, \Omega)$ is the projection $Q(-X^\alpha) = \delta^* \alpha$ onto $TS(M, \Omega)$ of the Hamiltonian vector field of $-\langle \Gamma(\nabla), \alpha \rangle$ on $\mathbb{T}(M, \Omega)$.

Recall from (4.13) the definition of the Poisson bracket $\{\cdot, \cdot\}$ on $(S(M, \Omega), \Omega)$.

**Lemma 5.3.** Let $(M, \Omega)$ be a compact symplectic manifold. For $\alpha, \beta \in S^2(M)$, the Poisson bracket $\{R_\alpha, R_\beta\}$ of $R_\alpha$ and $R_\beta$ is

$$\{R_\alpha, R_\beta\} = -\frac{2}{3}\langle \delta \alpha, \delta \beta \rangle + \frac{1}{3}\langle \alpha \ast \beta, C \rangle + \frac{1}{3}R(\alpha, \beta),$$

where $\alpha \ast \beta \in W^{2,2}(T^* M)$ is defined by $(\alpha \ast \beta)_{ijkl} = 2\alpha_{k[l[ij]} \beta_{j]l]} + 2\alpha_{l[ij} \beta_{k]} - (\beta \ast \alpha)_{ijkl}$. The Hamiltonian vector field on $(S(M, \Omega), \Omega)$ generated by $\{R_\alpha, R_\beta\}$ is $H_{\{R_\alpha, R_\beta\}} = -\frac{1}{3}(\delta^* (\alpha, \beta) + [\alpha, \beta]) = \frac{1}{3}(H_{R(\alpha, \beta)} - [\alpha, \beta]).$ (5.4)

**Proof.** Since $C_{ijkl} \in W^{2,2}(T^* M)$, the pairing (2.3) takes the form

$$\langle \alpha \ast \beta, C \rangle = 2 \int_M \alpha^{ik} \beta^{jk} R_{ijkl} \Omega_n.$$  

By (3.5) of Lemma 3.4,

$$\delta^* \beta_{ij} = -\frac{2}{3} \delta^* \beta_{ij} - \frac{2}{3} R^p (ij) \beta_{pq} + \frac{1}{3} (\beta, \text{Ric})_{ij}. \quad (5.5)$$

By definition of the Poisson bracket $\{R_\alpha, R_\beta\}$, Lemma 5.2, (5.5), and (3.4),

$$\{R_\alpha, R_\beta\} = \Omega(H_{R_\alpha}, H_{R_\beta}) = \langle \delta^* \alpha, \delta^* \beta \rangle = \langle \alpha, \delta^* \beta \rangle = -\frac{2}{3}\langle \delta \alpha, \delta \beta \rangle - \frac{2}{3} \int_M R^p (ij) \beta_{pq} \alpha^{ij} \Omega_n + \frac{1}{3}\langle \alpha, (\beta, \text{Ric}) \rangle$$

$$= -\frac{2}{3}\langle \delta \alpha, \delta \beta \rangle + \frac{1}{3}\langle \alpha \ast \beta, C \rangle + \frac{1}{3}\langle (\alpha, \beta), \text{Ric} \rangle.$$ (5.6)

By the second equality of (5.6) combined with (3.11), (3.4), and (2) of Lemma 3.5,

$$\Omega(\Pi, H_{\{R_\alpha, R_\beta\}}) = \delta_\Pi \{R_\alpha, R_\beta\} = \frac{1}{3}\langle (\alpha, \Pi), \delta^* \beta \rangle + \frac{1}{3}\langle \delta^* \alpha, (\beta, \Pi) \rangle$$

$$= \frac{1}{3}\langle (\alpha, \delta^* \beta) + (\delta^* \alpha, \beta), \Pi \rangle = -\frac{1}{3}\langle \Pi, \delta^* (\alpha, \beta) + [\alpha, \beta] \rangle,$$

which with Lemma 5.2 shows (5.4). ■
The identity (5.7) is essentially the same as (5.4). It yields a conceptual explanation of claim (2) of Lemma 3.5, that the Schouten bracket is a coboundary with respect to the algebraic bracket \((\cdot, \cdot)\), as it exhibits \([\cdot, \cdot]\) as the obstruction to making \((S^2(M), (\cdot, \cdot))\) act on an affine space. On the other hand, it means that \((S^2(M), (\cdot, \cdot))\) can be extended by \(S^3(M)\), using \([\cdot, \cdot]\) to produce a Lie algebra acting on \(S(M, \Omega)\). The idea developed in the remainder of this section and in Section 6 is that, although \(\mathcal{L}(\mathcal{F})\) does not act on \(S(M, \Omega)\), and so \(\Gamma\) (and correspondingly \(\text{Ric}\)) cannot be viewed as a moment map on \(S(M, \Omega)\), the action of \(\mathcal{L}(\mathcal{F})\) can be extended in such a way that a functional on \(S(M, \Omega)\) constructed from \(\text{Ric}\) yields a moment map for the action of the extended Lie algebra on \(S(M, \Omega)\), and the actions of \(\text{ham}(M, \Omega)\) and \(\mathcal{L}(\mathcal{F})\) can be combined and further extended in such a way that a functional on \(S(M, \Omega)\) constructed from \(\mathcal{X}\) and \(\text{Ric}\) yields a moment map for the action of the extended Lie algebra on \(S(M, \Omega)\).

As a vector space \(S^2(M) \oplus S^3(M)\) is identified with the quotient \(\mathcal{S}(S^2(M))/S^4(M)\). Since \(\mathcal{S}(S^2(M))/S^4(M)\) is an ideal of the Lie algebra \((\mathcal{S}(S^2(M)), [\cdot, \cdot])\) for the Lie bracket \([\cdot, \cdot]\) defined in (3.12), this bracket descends to a Lie bracket, also denoted \([\cdot, \cdot]\), on \((S^2(M) \oplus S^3(M))\), given explicitly by
\[
[a, b] = (\alpha_2, \beta_2) + [[\alpha_2, \beta_2] + (\alpha_2, \beta_3) - (\beta_2, \alpha_3)
\]
for \(\alpha, \beta \in S^2(M) \oplus S^3(M)\). For \(\alpha \in S^2(M) \oplus S^3(M)\) and \(\nabla \in S(M, \Omega)\), setting
\[
\alpha \cdot \nabla = \nabla + 3s^\alpha_2 + 3s_3,
\]
yields a symplectic affine action on \((S(M, \Omega))\) of the Lie algebra \((S^2(M) \oplus S^3(M), [\cdot, \cdot])\), where \([\cdot, \cdot]\) is the Lie bracket defined by (5.8). This is a special case of the more general Lemma 5.4.

**Lemma 5.4.** Let \((M, \Omega)\) be a symplectic manifold. Let \((\tilde{\mathcal{B}}_1(M), [\cdot, \cdot]_\tau)\) be the Lie algebra defined by (3.21) and (3.22). The map \(\pi^\tilde{B}_1: \tilde{\mathcal{B}}_1(M) \times S(M, \Omega) \to S(M, \Omega)\) defined by
\[
\pi^\tilde{B}_1(\nabla, \alpha) = \alpha \cdot \nabla + \mathcal{L}_\nabla(\alpha) + 3\delta^{\nabla}_\tau \alpha_2 + 3\alpha_3,
\]
for \((\nabla, \alpha) \in \tilde{\mathcal{B}}_1(M) \times S(M, \Omega)\), is a symplectic affine action of \((\tilde{\mathcal{B}}_1(M), [\cdot, \cdot]_\tau)\) on \((S(M, \Omega), \Omega)\). The stabilizer of \(\nabla \in S(M, \Omega)\) is the Lie subalgebra
\[
\tilde{\mathcal{B}}_1(M)_\nabla = \{ \alpha \in \tilde{\mathcal{B}}_1(M): 3\alpha_3 = -\mathcal{L}_\nabla(\alpha) - 3\delta^{\nabla}_\tau \alpha_2\}.
\]

**Proof.** Consider the more general map \(\tilde{\mathcal{B}}_1(M) \times S(M, \Omega) \to S(M, \Omega)\) defined by
\[
\alpha \cdot \nabla = \nabla + q\mathcal{L}_\nabla(\alpha_1) + r\delta^{\nabla}_\tau \alpha_2 + s\alpha_3,
\]
for \(q, r, s \in \mathbb{R}\). Let \(\rho: (\tilde{\mathcal{B}}_1(M), [\cdot, \cdot]_\tau) \to \text{aff}(S(M, \Omega))\) be the associated map given by \(\rho(\alpha)\nabla = \alpha \cdot \nabla - \nabla\). From the identities
\[
\delta^{\nabla}_\tau + \Pi \alpha_2 = \delta^{\nabla}_\tau \alpha_2 + \frac{1}{3}(\alpha_2, \Pi),
\]
\[
\mathcal{L}_\nabla(\alpha_1) = \mathcal{L}_{\alpha_1}^\tau (\nabla + \Pi) = \mathcal{L}_{\alpha_1}^\tau \nabla + \mathcal{L}_{\alpha_1}^\tau \Pi = \mathcal{L}_\nabla(\alpha_1) + [[\alpha_1, \Pi]],
\]

The flow of \(-R_\alpha\) has the form \(\nabla \to \nabla + s\delta^\alpha\). Lemma 5.3 can be viewed as showing that \(\pi(\alpha, \nabla) = \alpha \cdot \nabla = \nabla + s\delta^\alpha\) does not define an affine action of \((S^2(M), (\cdot, \cdot)) \simeq \mathcal{L}(\mathcal{F})\) on \(S(M, \Omega)\), and that the obstruction is given by the Schouten bracket. For \(\alpha, \beta \in S^2(M)\), using (3.9) and (3.11) yields
\[
\alpha \cdot (\beta \cdot \nabla) - \beta \cdot (\alpha \cdot \nabla) = \frac{1}{3}s^2 ((\alpha, \delta^\beta + (\delta^\alpha, \beta)) = \frac{1}{3}s^2 (\delta^\alpha (\alpha, \beta) + [[\alpha, \beta]])
\]
\[
= (\alpha, \beta) \cdot \nabla - \nabla + \frac{1}{3}s(s - 3)s^\alpha (\alpha, \beta) + \frac{1}{3}s^2 [\alpha, \beta].
\]
Hence the closest \(\alpha \cdot \nabla\) comes to defining an affine action is when \(s = 3\), in which case
\[
(\alpha \cdot (\beta \cdot \nabla) - \beta \cdot (\alpha \cdot \nabla)) - ((\alpha, \beta) \cdot \nabla - \nabla) = 3[[\alpha, \beta]].
\]
(5.7)

The identity (5.7) is essentially the same as (5.4). It yields a conceptual explanation of claim (2) of Lemma 3.5, that the Schouten bracket is a coboundary with respect to the algebraic bracket \((\cdot, \cdot)\), as it exhibits \([\cdot, \cdot]\) as the obstruction to making \((S^2(M), (\cdot, \cdot))\) act on an affine space.
(the first is (3.11)) it follows that the map $\alpha \to \mathcal{L}(\rho(\alpha))$ associating with $\alpha$ the linear endomorphism

$$
\mathcal{L}(\rho(\alpha))\Pi = q[\alpha_1, \Pi] + \frac{r}{3}(\alpha_2, \Pi),
$$
of $\mathcal{S}^3(M)$, satisfies $\rho(\alpha)(\nabla + \Pi) - \rho(\alpha)\nabla = \mathcal{L}(\rho(\alpha))\Pi$. This shows that $\rho(\alpha)$ is an affine transformation of $\mathcal{S}(M, \Omega)$ with linear part $\mathcal{L}(\rho(\alpha))$, and so, by Lemma 2.1, $\alpha$ is an affine action with linear part $\mathcal{L}(\alpha) = \text{Id} + \mathcal{L}(\rho(\alpha))$. By the symplectic $(M, \Omega)$-equivariance of $\Omega$ and (3.4), for $\Pi, \Upsilon \in T^* \mathcal{S}(M, \Omega)$,

$$
\Omega_{\rho(\alpha)\nabla}(\mathcal{L}(\rho(\alpha))\Pi, \Upsilon) + \Omega_{\rho(\alpha)\nabla}(\Pi, \mathcal{L}(\rho(\alpha))\Upsilon) = q (\langle [\alpha_1, \Pi], \Upsilon \rangle + \langle \Pi, [\alpha_1, \Upsilon] \rangle) + \frac{r}{3} (\langle [\alpha_2, \Pi], \Upsilon \rangle + \langle \Pi, [\alpha_2, \Upsilon] \rangle) = 0.
$$

This shows that $\rho(\alpha)$ is symplectic affine for all $\alpha \in \mathcal{B}_1(M)$. For $\alpha, \beta \in \mathcal{B}_1(M)$,

$$
\mathcal{L}(\rho(\alpha))\rho(\beta)\nabla = q[\alpha_1, \rho(\beta)\nabla] + \frac{r}{3}(\alpha_2, \rho(\beta)\nabla)
$$

$$
= q[\alpha_1, q\nabla\beta_1 + r\delta\nabla\beta_2 + s\beta_3] + \frac{r}{3}(\alpha_2, q\nabla\beta_1 + r\delta\nabla\beta_2 + s\beta_3)
$$

$$
= q^2[\alpha_1, L\nabla\beta_1] + qr([\alpha_1, \delta\nabla\beta_2] + \frac{1}{3}(\alpha_2, L\nabla\beta_1)) + q\tau[\alpha_1, \beta_3] + \frac{r^2}{3}(\alpha_2, \delta\nabla\beta_2) + \frac{rs}{3}(\alpha_2, \beta_3).
$$

Hence

$$
\mathcal{L}(\rho(\alpha))\rho(\beta)\nabla - \mathcal{L}(\rho(\beta))\rho(\alpha)\nabla = q^2([\alpha_1, L\nabla\beta_1] + [L\nabla\alpha_1, \beta_1]) + qr([\alpha_1, \delta\nabla\beta_2] + \frac{1}{3}(\alpha_2, L\nabla\beta_1) + \frac{1}{3}(\alpha_2, L\nabla\alpha_1) + \frac{1}{3}L\nabla\beta_1) + q\tau[\alpha_1, \beta_3] + \frac{r^2}{3}(\alpha_2, \delta\nabla\beta_2) + \frac{rs}{3}(\alpha_2, \beta_3).
$$

Since $d\alpha_1 = 0 = d\beta_1, \alpha^1_1, \beta^1_1 \in \text{symp}(M, \Omega)$, and

$$
[\alpha_1, L\nabla\beta_1] + [L\nabla\alpha_1, \beta_1] = [L_{\alpha^1_1}, L_{\beta^1_1}]\nabla = L\nabla([\alpha_1, \beta_1]).
$$

Using (5.13) and (3.9) to simplify (5.12) yields

$$
\mathcal{L}(\rho(\alpha))\rho(\beta)\nabla - \mathcal{L}(\rho(\beta))\rho(\alpha)\nabla = q^2L\nabla([\alpha_1, \beta_1]) + qr(L_{\alpha^1_1}\delta\nabla\beta_2 - \frac{1}{3}(\beta_2, L\nabla\alpha_1) - L_{\beta^1_1}\delta\nabla\alpha_2 + \frac{1}{3}(\alpha_2, L\nabla\beta_1)) + q\tau([\alpha_1, \beta_3] + [\alpha_3, \beta_1]) + \frac{r^2}{3}([\alpha_2, \beta_2] + \delta\nabla\alpha_2) + \frac{rs}{3}(\alpha_2, \beta_3) + (\alpha_3, \beta_2).
$$

By (3.11), there holds $\delta^*_{\phi_t}(\gamma) = \delta\nabla\nabla\gamma = \frac{1}{3}(\gamma, \phi_t\nabla - \nabla)$ for $\gamma \in \mathcal{S}^3(M)$ and the local flow $\phi_t$ generated by $Z \in \text{symp}(M, \Omega)$. Since $\delta^*_{\phi_t}(\gamma) = \phi_t^*(\delta\nabla\phi_{-t}^*(\gamma))$, differentiating at $t = 0$ yields

$$
[L_Z, \delta\nabla\gamma] = \left(\mathcal{L}_Z \delta\nabla\right)\gamma = \frac{1}{3}(\gamma, L(Z)\nabla).
$$

Applying (5.15) to simplify (5.14) yields

$$
\mathcal{L}(\rho(\alpha))\rho(\beta)\nabla - \mathcal{L}(\rho(\beta))\rho(\alpha)\nabla = q^2L\nabla([\alpha_1, \beta_1]) + qr\delta\nabla(q([\alpha_1, \beta_2] + [\alpha_2, \beta_1]) + \frac{r^2}{3}(\alpha_2, \beta_2)) + q\tau ([\alpha_1, \beta_3] + [\alpha_3, \beta_1]) + \frac{r^2}{3}([\alpha_2, \beta_2] + \delta\nabla(\alpha_2, \beta_2) + \frac{rs}{3}(\alpha_2, \beta_3) + (\alpha_3, \beta_2)).
$$

Hence

$$
(\mathcal{L}(\rho(\alpha))\rho(\beta) - \mathcal{L}(\rho(\beta))\rho(\alpha) - \rho([\alpha, \beta])\nabla)
$$

$$
= q(q - 1)L\nabla([\alpha_1, \beta_1]) + r\delta\nabla(q - 1)([\alpha_1, \beta_2] + [\alpha_2, \beta_1]) + \frac{r^2}{3}(\alpha_2, \beta_2)
$$

$$
+ (q - 1)s([\alpha_1, \beta_3] + [\alpha_3, \beta_1]) + \frac{r^2}{3}([\alpha_2, \beta_2] + \frac{s(q - 3)}{3}(\alpha_2, \beta_3) + (\alpha_3, \beta_2)).
$$

This vanishes if $q = 1$ and $r = 3 = s$. This yields (5.9) and shows that (5.9) defines a symplectic affine action. That the stabilizer of $\nabla \in \mathcal{S}(M, \Omega)$ has the form (5.10) is straightforward. ■
Corollary 5.5. Let \((M, \Omega)\) be a symplectic manifold. Let \((\hat{\mathfrak{h}}(M), [\cdot, \cdot])\) be the Lie algebra defined by \((3.12)\) and \((3.13)\). The map \(\pi^\hat{\mathfrak{h}} : \hat{\mathfrak{h}}(M) \times S(M, \Omega) \to S(M, \Omega)\) defined by

\[
\pi^\hat{\mathfrak{h}}(\alpha, \nabla) = \alpha \cdot \nabla = \nabla + \mathcal{L}_\nabla(\alpha_1) + 3\delta^* \alpha_2 + 3\alpha_3 = \nabla - \mathcal{K}_\nabla(\alpha_0) + 3\delta^* \alpha_2 + 3\alpha_3, \tag{5.16}
\]

for \(\alpha \in \hat{\mathfrak{h}}(M)\) and \(\nabla \in S(M, \Omega)\), is a symplectic affine action of \((\hat{\mathfrak{h}}(M), [\cdot, \cdot])\) on \((S(M, \Omega), \Omega)\). The stabilizer of \(\nabla \in S(M, \Omega)\) is the Lie subalgebra

\[
\hat{\mathfrak{h}}(M)\nabla = \{\alpha \in \hat{\mathfrak{h}}(M) : 3\alpha_3 = -\mathcal{L}_\nabla(\alpha_1) - 3\delta^* \alpha_2\}.
\]

Proof. By definition \(\pi^\hat{\mathfrak{h}} = \pi^\mathfrak{h} \circ (\pi_1 \times \text{Id}_{S(M, \nabla)})\) where \(\pi^\mathfrak{h}\) is defined by \((5.9)\). By Lemmas 3.11 and 5.4, \(\pi^\hat{\mathfrak{h}}\) is a symplectic affine action.

Remark 5.6. Consider the more general map \(\hat{\mathfrak{h}}(M) \times S(M, \Omega) \to S(M, \Omega)\) defined by

\[
\alpha \cdot \nabla = \nabla + p\mathcal{K}_\nabla(\alpha_0) + q\mathcal{L}_\nabla(\alpha_1) + r\delta^* \alpha_2 + sa_3, \tag{5.17}
\]

for \(p, q, r, s \in \mathbb{R}\). As in the proof of Lemma 5.4, it is straightforward to show that the associated map \(\rho : (\hat{\mathfrak{h}}(M), [\cdot, \cdot]) \to \text{diff}(S(M, \Omega))\) given by \(\rho(\alpha)\nabla = \alpha \cdot \nabla - \nabla\) is a symplectic affine transformation of \((S(M, \Omega), \Omega)\) with linear part

\[
\mathcal{L}(\rho(\alpha))\Pi = p[\delta^* \alpha_0, \Pi] + q[\alpha_1, \Pi] + r\delta^* (\alpha_2, \Pi) = (q-p)[\alpha_1, \Pi] + r\delta^* (\alpha_2, \Pi).
\]

Arguing as in the proof of Lemma 5.4 shows directly that \((5.17)\) defines a symplectic affine action if \(q = p + 1\) and \(r = 3 = s\). In this case, \((5.17)\) has the form \((5.16)\), whatever is the value of \(p\).

6 Hamiltonian action of the extended Lie algebra on \((S(M, \Omega), \Omega)\)

In this section \(M\) is supposed compact. This is needed only to guarantee convergence of the integrals that appear. With appropriate qualifications, everything extends to noncompact \(M\), but this requires fussing that would be distracting here.

6.1. A reference connection \(\nabla_0 \in S(M, \Omega)\) and a fixed \(\Pi \in T_{\nabla_0}S(M, \Omega)\) determine a function \(\Theta_{\nabla_0, \Pi}(\nabla) = \langle \Pi, \nabla - \nabla_0 \rangle\) on \(S(M, \Omega)\). Since \(\delta^* \Theta_{\nabla_0, \Pi}(\nabla) = \langle \Pi, \mathcal{K} \rangle\), the Hamiltonian vector field \(\mathcal{H}_{\Theta_{\nabla_0, \Pi}} = -\Pi\) is that generating the flow on \(S(M, \Omega)\) given by translation by \(-t\Pi\).

Given \(p, q, r \in \mathbb{R}\), \(\nabla_0 \in S(M, \Omega)\), and \(\alpha \in \hat{\mathfrak{h}}(M)\) define \(\mathcal{M}_{\alpha}^{p,q,r} : S(M, \Omega) \to \mathbb{R}\) by

\[
\mathcal{M}_{\alpha}^{p,q,r}(\nabla) = p\langle \alpha_0, \mathcal{K}(\nabla) \rangle + q\mathcal{R}\alpha_2(\nabla) + r\Theta_{\nabla_0, \alpha_3}(\nabla).
\]

The dependence of \(\mathcal{M}_{\alpha}^{p,q,r}\) on the reference connection \(\nabla_0\) is not indicated so that the notation does not become too cluttered. By Theorem 4.1, Lemma 5.2, and the preceding paragraph, the Hamiltonian vector field \(\mathcal{H}_{\mathcal{M}_{\alpha}^{p,q,r}}\) generated on \(S(M, \Omega)\) by \(\mathcal{M}_{\alpha}^{p,q,r}\) is

\[
\mathcal{H}_{\mathcal{M}_{\alpha}^{p,q,r}} = -p\mathcal{H}(\alpha_0) - q\delta^* \alpha_2 - r\alpha_3.
\]

In the case \(p = -1\), \(q = 3\), and \(r = 3\), there is written simply \(\mathcal{M}\) for \(\mathcal{M}^{1,3,3}\). Lemma 6.2 computes the Poisson brackets \(\{\mathcal{M}_\alpha, \mathcal{M}_\beta\}\). The following preliminary lemma is needed in its proof.
Lemma 6.1. Let \((M, \Omega)\) be a compact symplectic manifold. For \(\nabla = \nabla_0 + \Pi_{ij}^k, \nabla_0 \in \mathcal{S}(M, \Omega),\) \(\alpha, \beta \in \mathbb{S}^2(M),\) and \(\alpha \star \beta\) as defined in Lemma 5.3,
\[
\langle\alpha \star \beta, C(\nabla)\rangle - 2\langle\delta \nabla_0 \alpha, \delta \nabla_0 \beta\rangle = \langle\alpha \star \beta, C(\nabla_0)\rangle - 2\langle\delta \nabla_0 \alpha, \delta \nabla_0 \beta\rangle + \langle[\alpha, \beta], \Pi\rangle. \tag{6.1}
\]

Proof. For \(\alpha, \beta \in \mathbb{S}^2(M),\) \(C(\nabla) = C(\nabla_0) + d\nabla_0 \Pi_{ijkl} + 2\Pi_{pl} \Pi_{jkl} p,\) by (4.5), so
\[
\langle\alpha \star \beta, C(\nabla)\rangle - \langle\alpha \star \beta, C(\nabla_0)\rangle = \langle\alpha \star \beta, d\nabla_0 \Pi + \Pi \wedge \Pi\rangle.
\]
Integrating by parts and using \((\delta \nabla_0 \alpha, \Pi) = -3(\delta \nabla_0 \alpha)^p \Pi_{ijp}\) and (3.7) yields
\[
\langle\alpha \star \beta, d\nabla_0 \Pi\rangle = -\frac{2}{3} \langle\delta \nabla_0 \alpha, \Pi, \beta\rangle + \frac{2}{3} \langle(\delta \nabla_0 \beta, \Pi)\rangle + \langle[\alpha, \beta], \Pi\rangle. \tag{6.2}
\]
For \(\alpha, \beta \in \mathbb{S}^2(M),\)
\[
\alpha^{ab} \beta^{cd} \Pi_{acp} \Pi_{bdp} = \alpha^{ba} \beta_{dc} \Pi_{acp} \Pi_{bdp} = \alpha^{ab} \beta_{cd} \Pi_{bdp} \Pi_{ac} = -\alpha^{ab} \beta_{cd} \Pi_{acp} \Pi_{bdp}\tag{6.3}
\]
so that \(\alpha^{ab} \beta^{cd} \Pi_{acp} \Pi_{bdp} = 0.\) By (6.3),
\[
\langle\alpha \star \beta, \Pi \wedge \Pi\rangle = \frac{1}{2} \int_M \langle(\alpha \star \beta)_{ijkl} (\Pi \wedge \Pi)_{ijkl} \rangle \Omega_n = 2 \int \alpha^{ab} \Pi_{ab} \beta^{cd} \Pi_{cd} \Omega_n. \tag{6.4}
\]
From
\[
d\nabla_0 \alpha_{[ij]i_1...i_{k-1}} = d\nabla_0 \alpha_{[ij]i_1...i_{k-1}} - (k - 1) (\Pi_{i(i_1...i_{k-1})jp} - \Pi_{j(i_1...i_{k-1})ip})
\]
there follows
\[
\delta \nabla_0 \alpha_{i_1...i_{k-1}} = \delta \nabla_0 \alpha_{i_1...i_{k-1}} + (-1)^{k-1}(k - 1) \Pi^{pq} (\Pi_{i(p} \alpha_{i_1...i_{k-1})q)}\tag{6.5}
\]
Computing using (6.5), \((\delta \nabla_0 \alpha, \Pi) = -3(\delta \nabla_0 \alpha)^p \Pi_{ijp},\) and (6.4) yields
\[
\langle\delta \nabla_0 \alpha, \delta \nabla_0 \beta\rangle - \langle\delta \nabla_0 \alpha, \delta \nabla_0 \beta\rangle = -\frac{1}{3} \langle\delta \nabla_0 \alpha, \Pi, \beta\rangle + \frac{1}{3} \langle\alpha, \delta \nabla_0 \beta, \Pi\rangle + \frac{1}{2} \langle[\alpha, \beta], \Pi \wedge \Pi\rangle. \tag{6.6}
\]
Combining (6.2) and (6.6) yields (6.1).

Lemma 6.2. Let \((M, \Omega)\) be a compact symplectic manifold and fix \(\nabla_0 \in \mathcal{S}(M, \Omega).\) Define a skew-symmetric map \(\Sigma_{\nabla_0} : \tilde{\mathcal{H}}(M) \times \tilde{\mathcal{H}}(M) \rightarrow \mathbb{R}\) by
\[
\Sigma_{\nabla_0}(\alpha, \beta) = 3\langle(\alpha_2 \star \beta_2), C(\nabla_0)\rangle - 6\langle(\delta \nabla_0 \alpha_2, \delta \nabla_0 \beta_2)\rangle - 3\Omega(\Pi^{ab} \Omega_0, \beta_3) - 9\Omega(\alpha_3, \Omega^{ab} \beta_3) - 9\langle\alpha_3, \beta_3\rangle. \tag{6.7}
\]
For \(\alpha, \beta \in \tilde{\mathcal{H}}(M),\) the Poisson brackets on \((\mathcal{S}(M, \Omega), \Omega)\) of \(\mathcal{M_0}\) and \(\mathcal{M_\beta}\) satisfy
\[
\{\mathcal{M}_\alpha, \mathcal{M}_\beta\} - \mathcal{M}_{[\alpha, \beta]} = \Sigma_{\nabla_0}(\alpha, \beta),
\]
where \(\Sigma_{\nabla_0}(\alpha, \beta)\) is regarded as a constant function on \(\mathcal{S}(M, \Omega).\)

Proof. First, there are made some general observations needed later in the proof. Differentiating \(\mathcal{K}(\phi_t^* \nabla) = \mathcal{K}(\nabla) \circ \phi_t\) along the flow \(\phi_t\) of \(X \in \text{symp}(M, \Omega),\) and using (4.9) yields
\[
\mathcal{H}_\nabla \mathcal{H}_\nabla(f) = \{f, \mathcal{K}(\nabla)\}, \quad \angle(\mathcal{H}_\nabla(f), \mathcal{H}_\nabla(g)) = \langle\{f, g\}, \mathcal{K}(\nabla)\rangle = \langle\{f, g\}, \mathcal{K}(\nabla)\rangle. \tag{6.8}
\]
for \( f, g, \in C^\infty(M) \). Similarly, for \( f \in C^\infty(M) \) and \( \alpha \in S^2(M) \), by (4.2),

\[
\langle \mathcal{H}_\nabla(f), \delta^* \alpha \rangle = -\langle \delta^* \mathcal{H}_\nabla(f), \alpha \rangle = -\langle \mathcal{L}_\delta \text{Ric}(\nabla), \alpha \rangle
\]

\[
= \langle \text{Ric}(\nabla), \mathcal{L}_\delta \alpha \rangle = \langle \text{Ric}(\nabla), [\delta^* f, \alpha] \rangle = \mathcal{R}_{[\delta^* f, \alpha]}(\nabla). \tag{6.9}
\]

By definition of \( \mathcal{H}_\nabla \), for \( f \in C^\infty(M) \) and \( \alpha \in S^3(M) \),

\[
\langle \mathcal{H}_\nabla(f), \alpha \rangle = \langle \mathcal{H}_{\nabla_\alpha}(f), \alpha \rangle + \langle [\delta^* f, \alpha], \nabla - \nabla_0 \rangle. \tag{6.10}
\]

By (3.11), for \( \alpha_2, \beta_2 \in S^2(M) \) and \( \alpha_3, \beta_3 \in S^3(M) \),

\[
\langle \delta^*_\nabla \alpha_2, \beta_3 \rangle + \langle \alpha_3, \delta^*_\nabla \beta_2 \rangle
\]

\[
= \langle \delta^*_\nabla \alpha_2, \beta_3 \rangle + \langle \alpha_3, \delta^*_\nabla \beta_2 \rangle + \frac{1}{3} \langle (\alpha_2, \beta_3) + (\alpha_3, \beta_2), \nabla - \nabla_0 \rangle. \tag{6.11}
\]

For \( \alpha, \beta \in \hat{\text{H}}(M) \),

\[
\langle M^{p,q}_\alpha, M^{p,q}_\beta \rangle (\nabla) = \Omega(\hat{\text{H}}^{p,q}_\alpha, \hat{\text{H}}^{p,q}_\beta)
\]

\[
= p^2 \langle \mathcal{H}_\nabla(\alpha_0), \mathcal{H}_\nabla(\beta_0) \rangle + \frac{q^2}{3} \langle (\alpha_2 * \beta_2, C(\nabla)) - 2\langle \delta^* \mathcal{H}_\nabla(\alpha_0), \delta^* \mathcal{H}_\nabla(\beta_0) \rangle + \mathcal{R}_{[\alpha_2, \beta_2]}(\nabla) \rangle
\]

\[
+ r^2 \langle \langle \alpha_3, \beta_3 \rangle + \langle \alpha_3, \delta^* \nabla_0, [\alpha_2, \beta_2] \rangle \rangle - \langle \delta^* \mathcal{H}_\nabla(\alpha_0), \delta^* \mathcal{H}_\nabla(\beta_0) \rangle + \frac{q^2}{3} \langle \delta^* \nabla_0, [\alpha_2, \beta_2] \rangle \rangle
\]

\[
+ \langle \mathcal{H}_\nabla(\alpha_0), \beta_3 \rangle + \langle \alpha_3, \mathcal{H}_\nabla(\beta_0) \rangle \rangle + qr \langle \langle \delta^* \mathcal{H}_\nabla(\alpha_0), \beta_3 \rangle + \langle \alpha_3, \delta^* \nabla_0, \beta_2 \rangle \rangle. \tag{6.12}
\]

Simplifying (6.12) using (6.8), (6.9), (6.10), (6.1) of Lemma 6.1, (6.11), and Lemma 5.3 yields

\[
\langle M^{p,q}_\alpha, M^{p,q}_\beta \rangle (\nabla) = p^2 \langle \{\alpha_0, \beta_0\}, \mathcal{K}(\nabla) \rangle - \langle \mathcal{R}_{[\alpha_2, \beta_2]}(\nabla), \nabla \rangle + \frac{q^2}{3} \mathcal{R}_{[\alpha_2, \beta_2]}(\nabla)
\]

\[
- pr\langle \delta^* \nabla_0, [\alpha_2, \beta_2] \rangle + \langle \alpha_3, \beta_3 \rangle \rangle - \langle \delta^* \nabla_0, [\alpha_2, \beta_2] \rangle \rangle
\]

\[
+ \langle \mathcal{H}_\nabla(\alpha_0), \beta_3 \rangle + \langle \alpha_3, \mathcal{H}_\nabla(\beta_0) \rangle \rangle + qr \langle \langle \delta^* \mathcal{H}_\nabla(\alpha_0), \beta_3 \rangle + \langle \alpha_3, \delta^* \nabla_0, \beta_2 \rangle \rangle.
\]

If \( p = -1, q = 3, \) and \( r = 3 \), then, writing \( \mathcal{M}^{-1,3,3} = \mathcal{M} \),

\[
\langle \{ M_\alpha, M_\beta \} \rangle (\nabla) = \langle \{ \alpha_0, \beta_0 \}, \mathcal{K}(\nabla) \rangle + 3\mathcal{R}_{[\alpha_2, \beta_2]} + [\alpha_2, \beta_2] (\nabla)
\]

\[
+ 3\langle \delta^* \nabla_0, [\alpha_2, \beta_2] \rangle + \{ \alpha_2, \beta_2 \} (\nabla)
\]

\[
+ 9\langle \alpha_3, \beta_3 \rangle + 3\langle \delta^* \mathcal{H}_\nabla(\alpha_0), \delta^* \mathcal{H}_\nabla(\beta_0) \rangle - 2\langle \delta^* \mathcal{H}_\nabla(\alpha_0), \delta^* \mathcal{H}_\nabla(\beta_0) \rangle
\]

\[
- 3\langle \delta^* \mathcal{H}_\nabla(\alpha_0), \beta_3 \rangle + \langle \alpha_3, \mathcal{H}_\nabla(\beta_0) \rangle \rangle + 9\langle \delta^* \mathcal{H}_\nabla(\alpha_0), \beta_3 \rangle + \langle \alpha_3, \delta^* \nabla_0, \beta_2 \rangle \rangle
\]

\[
= \langle \mathcal{M}_{[\alpha, \beta]}(\nabla), \mathcal{M}_{[\alpha, \beta]}(\nabla) \rangle - \mathcal{M}_{[\alpha, \beta]}(\nabla). \tag{6.13}
\]

Since \( \mathcal{M} \) is not equivariant, the skew-symmetric map \( \Sigma_{\nabla_0} : \hat{\text{H}}(M) \times \hat{\text{H}}(M) \to \mathbb{R} \) is a 2-cocycle, the nonequivariance cocycle, for the Lie algebra cohomology of \( (\hat{\text{H}}(M), [\cdot, \cdot]) \) with coefficients in the trivial module \( \mathbb{R} \). The nonequivariance cocycles associated with different choices of reference connection \( \nabla_0 \in S(M, \Omega) \) are cohomologous. In what follows the dependence on the choice of \( \nabla_0 \) is sometimes suppressed, and there is written simply \( \Sigma \) instead of \( \Sigma_{\nabla_0} \). In a finite-dimensional setting it follows from the definition that the nonequivariance cocycle of a weakly Hamiltonian action is a constant function. In the infinite-dimensional setting here, the same conclusion would follow from the definition (6.7), had adequate analytic foundations been laid. Although they have not, Lemma 6.2 shows explicitly that \( \langle M_{[\alpha, \beta]}(\nabla) - M_{[\alpha, \beta]}(\nabla) \rangle = \mathcal{M}_{[\alpha, \beta]}(\nabla) \) is constant on \( S(M, \Omega) \).
6.2. Because \( \tilde{Z}_i(M) = \tilde{Z}(M) \cap \tilde{S}_i(M) = \tilde{S}_i(M) \) is an ideal in \( (\tilde{Z}_1(M), [\cdot, \cdot]) \) for any \( i \geq 2 \), the Lie bracket \([\cdot, \cdot]_\tau\) induces a Lie bracket (that will also be denoted \([\cdot, \cdot]_\tau\)) on the vector space \( \tilde{Z}_1(M) / \tilde{Z}_2(M) \). As a vector space \( \tilde{Z}_1(M) / \tilde{Z}_2(M) \) is identified with \( \text{symp}(M,\Omega)^7 \oplus \mathbb{S}^2(M) \oplus \mathbb{S}^3(M) \), and the Lie bracket is given explicitly by

\[
[\alpha, \beta]_\tau = [\alpha_1^\tau, \beta_1^\tau]^\gamma + (\alpha_2, \beta_2) - \mathcal{L}_{\beta_2^\tau} \alpha_2 + (\alpha_2, \beta_3) + (\alpha_3, \beta_2) - \mathcal{L}_{\beta_3} \alpha_2 - \mathcal{L}_{\beta_1^\tau} \alpha_3.
\]  

(6.13)

Although (6.13) satisfies the Jacobi identity by construction, this can also be checked by tedious computations using the compatibility and \( \text{symp}(M,\Omega) \)-equivariance of \([\cdot, \cdot]_\tau\).

Since \( \tilde{B}_4(M) = \tilde{B}(M) \cap \tilde{S}_4(M) \) satisfies \( \tilde{B}_1(M) = \tilde{Z}_4(M) \) and \( \tilde{B}_1(M) \) is a subalgebra of \( (\tilde{Z}_1(M), [\cdot, \cdot]_\tau) \), \( (\tilde{B}_1(M) / \tilde{B}_4(M), [\cdot, \cdot]) \) is a Lie algebras of \( (\tilde{Z}_1(M) / \tilde{Z}_4(M), [\cdot, \cdot]_\tau) \) and the exact sequence (3.23) of Lemma 3.11 descends to give an exact sequence

\[
\{0\} \longrightarrow \mathbb{R} \xrightarrow{\iota} (\tilde{H}(M) / \tilde{H}_4(M), [\cdot, \cdot]) \xrightarrow{\pi_1} (\tilde{B}_1(M) / \tilde{B}_4(M), [\cdot, \cdot]_\tau) \longrightarrow \{0\}
\]

doing Lie algebras, where \( \iota \) is the restriction to the constant functions \( \mathbb{R} \subset C^\infty(M) \) of the map \( \iota: C^\infty(M) \rightarrow \tilde{H}(M) \) defined by \( \iota(f) = -f + \delta^* f \) (see Lemma 3.7), and embeds \( \mathbb{R} \) in \( \tilde{H}(M) = C^\infty(M) \) as the constant functions (but with a sign change). The bracket \([\cdot, \cdot]_\tau\) descends to a Lie bracket, also denoted \([\cdot, \cdot]_\tau\) on the quotient \( \tilde{H}(M)(\iota(\mathbb{R}) \oplus \tilde{H}_4(M)) \) by the Lie ideal \( \iota(\mathbb{R}) \oplus \tilde{H}_4(M) \).

For a compact symplectic manifold \( (M,\Omega) \), the linear map \( c: (\tilde{H}(M), [\cdot, \cdot]) \rightarrow \mathbb{R} \) defined by

\[
c(\alpha) = \text{vol}_{\Omega_n}(M)^{-1} \int_M \alpha_0 \Omega_n = \text{vol}_{\Omega_n}(M)^{-1}\langle \alpha_0, 1 \rangle,
\]

is a Lie algebra homomorphism because, by (3.17), \([\alpha, \beta]_\tau = -\{\alpha_0, \beta_0\}\), and the integral of a Poisson bracket vanishes. Define \( \nu: \tilde{H}(M) \rightarrow \tilde{H}(M) \) by \( \nu(\alpha) = \alpha + \iota(c(\alpha)) \) (so \( \nu(\alpha_0) = \alpha_0 - c(\alpha_0) \)). Because \( \iota(\mathbb{R}) \) is central in \( (\tilde{H}(M), [\cdot, \cdot]) \), \( \nu \) is a Lie algebra homomorphism of \( (\tilde{H}(M), [\cdot, \cdot]) \). Because \( \nu \circ \nu = \nu, \pi_1 \circ \nu = \pi_1 \), and ker \( \nu = \iota(\mathbb{R}) \), it follows that the maps \( \nu \) and \( \pi_1 \) induce Lie algebra isomorphisms

\[
(\tilde{H}(M)(\iota(\mathbb{R}) \oplus \tilde{H}_4(M)), [\cdot, \cdot]) \cong (\nu(\tilde{H}(M))(\tilde{H}_4(M)), [\cdot, \cdot]) \cong (\tilde{B}_1(M) / \tilde{B}_4(M), [\cdot, \cdot]_\tau).
\]

When \( M \) is compact, the Lie algebra \( \text{ham}(M,\Omega) \) is isomorphic with the mean zero elements of \( C^\infty(M) \) equipped with the Poisson bracket. It follows from Lemma 3.7 that the map \( \tilde{H}(M) \rightarrow \mathfrak{h} = \text{ham}(M,\Omega) \oplus \mathbb{S}^2(M) \oplus \mathbb{S}^3(M) \) defined by \( \alpha \mapsto -\alpha_0 + \alpha_2 + \alpha_3 \) induces a linear isomorphism \( \tilde{H}(M)(\iota(\mathbb{R}) \oplus \tilde{H}_4(M)) \rightarrow \mathfrak{h} \) with inverse \( j: \mathfrak{h} \rightarrow \tilde{H}(M)(\iota(\mathbb{R}) \oplus \tilde{H}_4(M)) \) given by \( j(f + \alpha_2 + \alpha_3) = \iota(f) + \alpha_2 + \alpha_3 = -f + \delta^* f + \alpha_2 + \alpha_3 \in \tilde{H}(M)(\iota(\mathbb{R}) \oplus \tilde{H}_4(M)) \), and such that the pullback via \( j \) of the Lie bracket \([\cdot, \cdot]_\tau\), also to be written \([\cdot, \cdot]_\tau\), agrees, when restricted to the subspace \( \text{ham}(M,\Omega) \), with the Poisson bracket.

Since for the symplectic affine action (5.16) the action of \( \mathbb{R} \oplus \tilde{H}_4(M) \) is trivial, the action (5.16) descends to a symplectic affine action of \( (\tilde{H}(M)(\iota(\mathbb{R}) \oplus \tilde{H}_4(M)), [\cdot, \cdot]) \) on \( \mathcal{S}(M,\Omega) \). Composing with \( j \) yields a symplectic affine action of \( (\tilde{h}, [\cdot, \cdot]) \) on \( \mathcal{S}(M,\Omega) \). Precisely, the symplectic affine action \( \pi^h_\tau \) of \( \alpha \in \mathfrak{h} \) on \( \mathcal{S}(M,\Omega) \) is defined to be \( \pi^h_\tau(\alpha, \nabla) = \pi^H_\tau(j(\alpha), \nabla) \); this makes sense because \( \tilde{H}_4(M) \) acts trivially on \( \mathcal{S}(M,\Omega) \). Lemma 6.3 summarizes the preceding as showing that this action is weakly Hamiltonian. In the statement, the dual \( \mathfrak{h}^* \) is identified with \( \mathfrak{h} \) via the pairing \( \langle \cdot, \cdot \rangle \). Precisely, the pairing \( \langle \cdot, \cdot \rangle \) extends to elements \( \alpha \) and \( \beta \) of \( \tilde{S}(M) \) (or of \( \tilde{H}(M) \) or \( \mathfrak{h} \)) by linearity as \( \langle \alpha, \beta \rangle = \sum_{k \geq 0} \langle \alpha_k, \beta_k \rangle \). Note that, if \( M \) is compact and \( \alpha, \beta \in \tilde{H}(M) \), then \( \langle \alpha_1, \beta_1 \rangle = \langle \delta^* \alpha_0, \delta^* \beta_0 \rangle = \int_M \langle \alpha_0, \beta_0 \rangle \Omega_n = 0 \). Since \( M \) is compact elements of the linear duals of any of \( \tilde{H}(M) \), \( \mathfrak{h} \), etc. can be identified with elements of the space itself via the pairing \( \langle \cdot, \cdot \rangle \).
Lemma 6.3. Let \((M, \Omega)\) be a compact symplectic manifold. The action \(\pi^h\) of \((\mathfrak{h}, [\cdot, \cdot])\), on \(\mathcal{S}(M, \Omega)\) induced by the symplectic affine action (5.16) is weakly Hamiltonian with (nonequivalent) moment map \(\hat{\mathcal{M}} : \mathcal{S}(M, \Omega) \to \mathfrak{h}^*\) given by

\[
\langle \hat{\mathcal{M}}(\nabla), \alpha \rangle = \mathcal{M}_{\hat{\mathfrak{h}}(\alpha)}(\nabla) = \langle \alpha_0, \mathcal{X}(\nabla) \rangle + 3\mathfrak{X}_\alpha(\nabla) + 3\mathcal{G}_{\nabla_0, \alpha}(\nabla),
\]

for any fixed \(\nabla_0 \in \mathcal{S}(M, \Omega)\) and any \(\alpha = \alpha_0 + \alpha_2 + \alpha_3 \in \mathfrak{h}\).

**Proof.** By definition of the symplectic affine action \(\pi_{\hat{\mathfrak{h}}}\), for \(\alpha \in \hat{H}(M)\),

\[
\hat{H}_{\mathfrak{H}_\alpha}(\nabla) = \mathfrak{H}(\alpha_0) - 3\delta_\alpha \alpha_2 - 3\alpha_3 + \frac{d}{dt}|_{t=0} \pi_t\hat{\mathfrak{h}}(-t\alpha, \nabla).
\]

This shows that \(\hat{H}_{\mathfrak{H}_\alpha}\) is the vector field generated by the action of \(\alpha\) on \(\mathcal{S}(M, \Omega)\), so \(\mathcal{M}(\nabla)\) is a (nonequivalent) moment map for the action of \(\hat{H}(M)\) on \(\mathcal{S}(M, \Omega)\), and this action is weakly Hamiltonian. The claim follows formally from the definitions of \((\mathfrak{h}, [\cdot, \cdot])\) and \(\pi^h\). 

---

Fix \(\nabla_0 \in \mathcal{S}(M, \Omega)\). Define \(\hat{\mathfrak{h}} = C^\infty(M) \oplus \mathfrak{S}^2(M) \oplus \mathfrak{S}^3(M)\) and extend the pairing \(\langle \cdot, \cdot \rangle\) to \(\hat{\mathfrak{h}}\) as for \(\mathfrak{h}\). The linear isomorphism \(j : \mathfrak{h} \to \hat{H}(M)/(\iota(\mathbb{R}) \oplus \mathfrak{H}_4(M))\) lifts to a linear isomorphism \(\hat{j} : \hat{\mathfrak{h}} \to \hat{H}(M)/(\hat{\mathfrak{H}}_4(M)\) (defined in the same way as \(j\), by \(\hat{j}(f + \alpha_2 + \alpha_3) = \iota(f) + \alpha_2 + \alpha_3 \in \hat{H}(M)/(\hat{\mathfrak{H}}_4(M)\)). As for \(j\), the pullback to \(\hat{\mathfrak{h}}\) via \(\hat{j}\) of the Lie bracket \([\cdot, \cdot]\), again also to be written \([\cdot, \cdot]\), agrees, when restricted to the subspace \(C^\infty(M)\), with the Poisson bracket. Since \(\Sigma\) vanishes when restricted to \(\hat{\mathfrak{H}}_4(M)\) in either argument, it can be viewed as a pairing on \(\hat{H}(M)/(\hat{\mathfrak{H}}_4(M)\), so it makes sense to define a Lie bracket \(\langle \cdot, \cdot \rangle\) on \(\hat{\mathfrak{h}}\) by

\[
\langle (\alpha, \beta) \rangle = [\alpha, \beta] + \Sigma(\hat{j}(\alpha), \hat{j}(\beta)),
\]

(6.14)

Here \(\Sigma\) is regarded as taking values in the subspace of constant functions \(\mathbb{R} \subset C^\infty(M)\). Since \(\iota(\mathbb{R}) = \ker \nu\), \(\nu\) descends from \(\hat{H}(M)\) to give a Lie algebra homomorphism

\[
\nu : (\hat{H}(M)/(\hat{\mathfrak{H}}_4(M), [\cdot, \cdot])) \to (\hat{H}(M)/(\iota(\mathbb{R}) \oplus \mathfrak{H}_4(M), [\cdot, \cdot])).
\]

The linear map \(\hat{\nu} : \hat{j} \to j\) defined by \(\hat{j}(f + \alpha_2 + \alpha_3) = f - c(f) + \alpha_2 + \alpha_3\) satisfies \(j \circ \hat{\nu} = \nu \circ \hat{j}\), so is a Lie algebra homomorphism \(\hat{\nu} : (\hat{\mathfrak{h}}, \langle \cdot, \cdot \rangle) \to (\mathfrak{h}, [\cdot, \cdot])\). In what follows it is convenient, and should cause no confusion, to omit the bar and write simply \(\nu\) in place of \(\hat{\nu}\). The Lie algebra \((\mathfrak{h}, \langle \cdot, \cdot \rangle)\) is the central extension of \((\mathfrak{h}, [\cdot, \cdot])\) corresponding to the 2-cocycle \(\Sigma\). Precisely, the following sequence of Lie algebras is exact:

\[
\{0\} \to \mathbb{R} \to \langle \mathfrak{h}, \langle \cdot, \cdot \rangle \rangle \stackrel{\nu}{\longrightarrow} (\mathfrak{h}, [\cdot, \cdot]) \longrightarrow \{0\}.
\]

The action \(\hat{\pi}^h(\alpha, \nabla)\) of \((\hat{\mathfrak{h}}, \langle \cdot, \cdot \rangle)\) on \(\mathcal{S}(M, \Omega)\) is defined to be the action \(\pi^h(\nu(\alpha), \nabla)\) of its image via the projection \(\nu\), so that the center \(\mathbb{R}\) of \((\hat{\mathfrak{h}}, \langle \cdot, \cdot \rangle)\) acts trivially.

**Theorem 6.4.** On a compact symplectic manifold \((M, \Omega)\), for any \(\nabla_0 \in \mathcal{S}(M, \Omega)\) the action of \((\hat{\mathfrak{h}}, \langle \cdot, \cdot \rangle)\) on \((\mathcal{S}(M, \Omega), \Omega)\) is Hamiltonian, with equivariant moment map \(\mathcal{M}^* : \mathcal{S}(M, \Omega) \to \hat{\mathfrak{h}}^*\) given by

\[
\langle \mathcal{M}^*(\nabla), \alpha \rangle = \langle \mathcal{M}(\nabla), \nu(\alpha) \rangle - c(\hat{j}(\alpha)) = \mathcal{M}_{\hat{\mathfrak{h}}(\alpha)}(\nabla) - c(\hat{j}(\alpha)) = \mathcal{M}_{\nu(\alpha)}(\nabla) + c(\alpha_0)
\]

for any \(\alpha = \alpha_0 + \alpha_2 + \alpha_3 \in \hat{\mathfrak{h}}\), where \(\mathcal{M}\) is determined by \(\nabla_0\) as in Lemma 6.3.

**Proof.** Let \(\alpha, \beta \in \hat{\mathfrak{h}}\). It is convenient to write \(\hat{\pi}_\alpha\) and \(\mathcal{M}(\alpha)\) in place of \(\langle \hat{\pi}(\nabla), \alpha \rangle\) and \(\langle \mathcal{M}(\nabla), \nu(\alpha) \rangle\).
By (4.4), when $M$ is compact, the constant $\kappa(M, \Omega) = \int_M \mathcal{K}(\nabla) \Omega$ does not depend on the choice of $\nabla \in \mathcal{S}(M, \Omega)$. For $\beta \in \hat{H}(M)$,

$$\mathcal{M}_{\nu(\beta)}(\nabla) - \mathcal{M}_{\beta}(\nabla) = \langle c(\beta), \mathcal{K}(\nabla) \rangle = c(\beta)\kappa(M, \Omega)$$

is a constant function on $\mathcal{S}(M, \Omega)$. It follows that $\hat{M}_\alpha$, $\mathcal{M}_{\nu(\alpha)}$, and $\mathcal{M}_{\nu(\beta)}$ differ by constant functions, so they generate on $\mathcal{S}(M, \Omega)$ the same Hamiltonian vector field. This is the vector field generated by the symplectic affine action of $\nu(\alpha)$ on $\mathcal{S}(M, \Omega)$, and this equals the lifted action of $\alpha$ on $\mathcal{S}(M, \Omega)$. This suffices to show that $\hat{H}$ is a moment map for the action of $(\mathfrak{h}, (\cdot, \cdot))$ on $\mathcal{S}(M, \Omega)$, provided the equivariance of $\hat{H}$ can be shown.

For $\alpha, \beta \in \hat{H}(M)$, since, in the definition (6.7) of $\Sigma(\alpha, \beta)$, there enter $\alpha_0$ and $\beta_0$ only via $H_{\alpha_0}$ and $H_{\beta_0}$, and, by the preceding remarks, these vector fields equal $H_{\mathcal{M}_{\nu(\alpha)}}$ and $H_{\mathcal{M}_{\nu(\beta)}}$, there holds $\Sigma(\nu(\alpha), \nu(\beta)) = \Sigma(\alpha, \beta)$.

For $\alpha, \beta \in \mathfrak{h}$, because $[H(M), H(M)] \subset \ker c$,

$$c(\hat{\mathcal{J}}(\alpha, \beta)) = c(\hat{\mathcal{J}}(\alpha, \beta)) + c(\hat{\mathcal{J}}(\Sigma(\hat{\mathcal{J}}(\alpha), \hat{\mathcal{J}}(\beta))))$$

$$= c(\hat{\mathcal{J}}(\alpha, \beta)) + \hat{\mathcal{J}}(\hat{\mathcal{J}}(\alpha), \hat{\mathcal{J}}(\beta)) = -\hat{\mathcal{J}}(\hat{\mathcal{J}}(\alpha), \hat{\mathcal{J}}(\beta)). \quad (6.15)$$

By the definitions of $\hat{H}$ and $\Sigma$, for $\alpha, \beta \in \mathfrak{h}$,

$$\{\hat{H}_\alpha, \hat{H}_\beta\} = \{\mathcal{M}_{\nu(\alpha)}, \mathcal{M}_{\nu(\beta)}\} = \{\mathcal{M}_{\nu(\alpha), \nu(\beta)}\}$$

$$= \mathcal{M}_{\mathcal{M}_{\nu(\alpha), \nu(\beta})} + \mathcal{M}(\mathcal{J}(\nu(\alpha), \mathcal{J}(\nu(\beta))) = \mathcal{M}(\mathcal{J}(\nu(\alpha), \nu(\beta))) + \mathcal{J}(\mathcal{J}(\nu(\alpha), \nu(\beta)))$$

$$= \mathcal{M}_{\mathcal{M}_{\nu(\alpha), \nu(\beta})} + \mathcal{M}(\mathcal{J}(\nu(\alpha), \mathcal{J}(\nu(\beta)))) = \mathcal{M}_{\mathcal{M}_{\nu(\alpha), \nu(\beta})} = \mathcal{M}(\mathcal{J}(\nu(\alpha), \mathcal{J}(\nu(\beta))))$$

where the penultimate equality follows from (6.15). ■

Viewed as a map taking values in the dual of $\hat{H}(M)/(\iota(\mathbb{R}) \oplus \hat{H}_4(M))$, the map $\hat{H}(\nabla)$ is identified via $\langle \cdot, \cdot \rangle$ with the element

$$j(\mathcal{K}(\nabla) + 3\text{Ric}(\nabla) + 3(\nabla - \nabla_0)) = -\mathcal{K}(\nabla) + \delta^*\mathcal{K}(\nabla) + 3(\text{Ric}(\nabla) + 3(\nabla - \nabla_0))$$

of $\hat{H}(M)/(\iota(\mathbb{R}) \oplus \hat{H}_4(M))$, so the map $\hat{H}(\nabla) \in \mathfrak{h}^*$ is identified via $\langle \cdot, \cdot \rangle$ with

$$\mathcal{K}(\nabla) + 3\text{Ric}(\nabla) + 3(\nabla - \nabla_0) \in \mathfrak{h}.$$ 

Similarly $\hat{H}(\nabla) \in \mathfrak{h}^*$ is identified with

$$\hat{H}(\nabla) = \text{vol}_{\mathcal{O}_n}(M)^{-1} + \mathcal{K}(\nabla) + 3\text{Ric}(\nabla) + 3(\nabla - \nabla_0) \in \mathfrak{h}, \quad (6.16)$$

where $\mathfrak{h}^*$ is identified with $\hat{h}$ using $\langle \cdot, \cdot \rangle$, $\mathcal{K}(\nabla) = \mathcal{K}(\nabla) - \kappa(M, \Omega)$ is the mean zero part of $\mathcal{K}(\nabla)$, and the constant function $\text{vol}_{\mathcal{O}_n}(M)^{-1}$ is identified via $\langle \cdot, \cdot \rangle$ with the map $f \in C^\infty(M) \rightarrow c(f)$.

On a compact symplectic manifold fix $\nabla_0 \in \mathcal{S}(M, \Omega)$ with corresponding moment map $\hat{H}: \mathcal{S}(M, \Omega) \rightarrow \mathfrak{h}^*$ and define $\mathcal{N}: \mathcal{S}(M, \Omega) \rightarrow \mathbb{R}$ by

$$\mathcal{N}(\nabla) = \langle \hat{H}(\nabla), \hat{H}(\nabla) \rangle = \text{vol}_{\mathcal{O}_n}(M)^{-1}(1 - \kappa(M, \Omega)^2) + \mathcal{E}(\nabla) + 18\mathcal{R}(\nabla),$$

where $\mathcal{E}(\nabla) = \int M \mathcal{K}(\nabla)^2 \mathcal{O}_n$ and $\mathcal{R}(1)$ is defined in (4.11). By (4.10) and (4.12), the first variation of $\mathcal{N}$ is

$$\delta_{\Pi} \mathcal{N}(\nabla) = 2\mathcal{O}_n \mathcal{K}(\nabla) + 9\delta^*\text{Ric},$$

so the critical points of $\mathcal{N}$ are the solutions of $\mathcal{K}(\nabla) + 9\delta^*\text{Ric} = 0.$
6.3. For $s, t \in \mathbb{R}^\times$ let $\Psi_{s,t} \in \text{End}(\tilde{S}(M))$ be the linear automorphism defined in Lemma 3.6. Note that, although $\Psi_{s,t}$ does preserve $B_1(M)$, it does not preserve $\tilde{H}(M)$. Rather, the subspace $\tilde{H}^{s,t}(M) = \{ \alpha \in \tilde{H}(M); s\alpha_1 = -t\delta^*\alpha_0 \}$ is a subalgebra of $(\tilde{S}(M), [\cdot, \cdot]_{s,t})$ and $\Psi_{s,t}$ maps it isomorphically onto $\tilde{H}(M)$. Hence, composing $\tilde{\pi}_h$ with $\Psi_{s,t} \times \text{Id}_{\tilde{S}(M, \Omega)}$ yields a degenerated symplectic affine action of $(\tilde{\tilde{H}}^{s,t}(M), [\cdot, \cdot]_{s,t})$ on $S(M, \Omega)$ given by

$$\tilde{\pi}_{s,t}(\alpha, \nabla) = \nabla + t\mathcal{L}_\nabla(\alpha_1) + 3s\delta^*\alpha_2 + 3t^{-1}s^2\alpha_3$$

$$= \nabla - s^{-1}t^2\mathcal{L}_\nabla(\alpha_0) + 3s\delta^*\alpha_2 + 3t^{-1}s^2\alpha_3.$$ 

These actions are all isomorphic as long as $s$ and $t$ are both nonzero. The extremal case $(s, t) = (0, 1)$ yields a symplectic affine action of $(\tilde{\tilde{H}}(M), [\cdot, \cdot]_1)$ on $S(M, \Omega)$ that corresponds up to a constant factor with the usual action of $\text{ham}(M, \Omega)$ on $S(M, \Omega)$. All the constructions made starting from $(\tilde{\tilde{H}}(M), [\cdot, \cdot])$ can be repeated with $(\tilde{\tilde{H}}^{s,t}(M), [\cdot, \cdot]_{s,t})$ in its place (this requires modifying the maps $j$ and $j$). There results a moment map $\hat{\pi}^{s,t}$ for a Lie bracket $([\cdot, \cdot])_{s,t}$ on $\tilde{h}$.

However, it is easier, and equivalent, to introduce the parameters $(s, t)$ at the end of the construction, by viewing $\Psi_{s,t}$ as a linear endomorphism of $\tilde{h}$, and defining $(([\cdot, \cdot])_{s,t}$ directly to be the pullback of $([\cdot, \cdot])_{s,t}$ via $\Psi_{s,t}$. Here there has to be made precise how to view $\Psi_{s,t}$ as an endomorphism of $\tilde{h}$. It is simpler still to work with a modified map $\psi_{s,t} : \tilde{h} \to \tilde{h}$ defined by $\psi_{s,t}(\alpha_0 + \alpha_2 + \alpha_3) = t\alpha_0 + s\alpha_2 + t^{-1}s^2\alpha_3$, and to define $(([\cdot, \cdot])_{s,t}$ to be the pullback of $([\cdot, \cdot])_{s,t}$ via $\psi_{s,t}$. The composed action $\hat{\pi}_h \circ (\psi_{s,t} \times \text{Id}_{\tilde{h}})$ of $(\tilde{h}, ([\cdot, \cdot])_{s,t})$ is symplectic affine, and the corresponding moment map $\hat{\pi}^{s,t} : S(M, \Omega) \to \tilde{h}^*$ is $\hat{\pi}^{s,t} = \psi_{s,t}^* \circ \hat{\pi}_h$ where $\psi_{s,t}^* \in \text{End}(\tilde{h})$ is the linear dual of $\psi_{s,t}$. The corresponding quadratic functional $N_{s,t}(\nabla)$ is

$$N_{s,t}(\nabla) = \langle \hat{\pi}^{s,t}(\nabla), \hat{\pi}^{s,t}(\nabla) \rangle$$

$$= t^2\mathcal{V}(\nabla, \nabla) - t^2\kappa(M, \Omega)^2 + t^2\mathcal{E}(\nabla) + 18s^2\mathcal{R}(\nabla) + 18s^2\mathcal{R}(\nabla) + 18s^2\mathcal{R}(\nabla).$$

By (4.10) and (4.12), the first variation of $N_{s,t}$ is

$$\delta N_{s,t}(\nabla) = 2\nabla \left( t^2\mathcal{H}(\nabla) + s^2\delta^*\text{Ric}, \Pi \right),$$

so the critical points of $N_{s,t}$ are the solutions of (1.2).

Acknowledgements I thank the anonymous referees for their thoughtful criticisms and detailed corrections which helped improve the article, particularly the exposition in Section 6.

References

[1] Atiyah M.F., Bott R., The Yang–Mills equations over Riemann surfaces, *Philos. Trans. Roy. Soc. London Ser. A* **308** (1983), 523–615.

[2] Bieliavsky P., Cahen M., Gutt S., Rawnsley J., Schwachhöfer L., Symplectic connections, *Int. J. Geom. Methods Mod. Phys.* **3** (2006), 375–420, math.SG/0511194.

[3] Bolsinov A.V., Compatible Poisson brackets on Lie algebras and the completeness of families of functions in involution, *Math. USSR Izv.* **38** (1992), 69–90.

[4] Bolsinov A.V., Borisov A.V., Compatible Poisson brackets on Lie algebras, *Math. Notes* **72** (2002), 10–30.

[5] Bourgeois F., Cahen M., Can one define a preferred symplectic connection?, *Rep. Math. Phys.* **43** (1999), 35–42.

[6] Bourgeois F., Cahen M., A variational principle for symplectic connections, *J. Geom. Phys.* **30** (1999), 233–265.

[7] Cahen M., Gutt S., Moment map for the space of symplectic connections, in Liber Amicorum Delanghe, Editors F. Brackx, H. De Schepper, Gent Academia Press, 2005, 27–36.
[8] Fedosov B., Deformation quantization and index theory, *Mathematical Topics*, Vol. 9, Akademie Verlag, Berlin, 1996.

[9] Fedosov B., The Atiyah–Bott–Patodi method in deformation quantization, *Comm. Math. Phys.* **209** (2000), 691–728.

[10] Fox D.J.F., Critical symplectic connections on surfaces, arXiv:1410.1468.

[11] Golubchik I.Z., Sokolov V.V., Compatible Lie brackets and integrable equations of the principal chiral field model type, *Funct. Anal. Appl.* **36** (2002), 172–181.

[12] Guillemin V., The integrability problem for $G$-structures, *Trans. Amer. Math. Soc.* **116** (1965), 544–560.

[13] Gutt S., Remarks on symplectic connections, *Lett. Math. Phys.* **78** (2006), 307–328.

[14] Kobayashi S., Transformation groups in differential geometry, *Ergebnisse der Mathematik und ihrer Grenzgebiete*, Vol. 70, Springer-Verlag, Berlin – Heidelberg, 1995.

[15] Kostant B., The Weyl algebra and the structure of all Lie superalgebras of Riemannian type, *Transform. Groups* **6** (2001), 215–226, math.RT/0106171.

[16] McDuff D., Salamon D., Introduction to symplectic topology, Oxford Mathematical Monographs, The Clarendon Press, Oxford University Press, New York, 1995.

[17] Penrose R., Rindler W., Spinors and space-time, Vol. 1. Two-spinor calculus and relativistic fields, *Cambridge Monographs on Mathematical Physics*, Cambridge University Press, Cambridge, 1984.

[18] Sternberg S., Lectures on differential geometry, Prentice-Hall, Inc., Englewood Cliffs, N.J., 1964.

[19] Tamarkin D.E., Topological invariants of connections on symplectic manifolds, *Funct. Anal. Appl.* **29** (1995), 258–267.

[20] Wald R.M., General relativity, *University of Chicago Press*, Chicago, IL, 1984.