ENERGY CONTENT OF GRAVITATION AS A WAY TO QUANTIFY BOTH ENTROPY AND INFORMATION GENERATION IN THE EARLY UNIVERSE

ANDREW WALCOTT BECKWITH
beckwith@aibep.org, abeckwith@uh.edu

American Institute of Beam Energy Propulsion, life member

Based upon Alcubierre’s formalism about energy flux of gravitational waves, as well as Saundër’s treatment of temperature dependence of the Hubble parameter in the early universe, we initiate a particle count treatment of gravitons, and subsequent entropy generation which gives, via the standard model treatment of the FRW metric, a way to explain/justify a value of entropy of the order of $S \sim 10^6 - 10^7$ at the very onset of inflation.

1 Introduction

The supposition advanced in this article is that relic energy flux initially is central to making predictions as to $S_{\text{entropy}} \sim n_f$, where $n_f$ is a ‘particle count’ per phase space ‘volume’ in the beginning of inflation. The author is aware of how many researchers have linked relic GW as to initial phase transitions as to the electro weak phase transition, in cosmology. The supposition is that $S \sim 10^6 - 10^7$ is demonstrable as an initial entropy/information count in the onset of inflation if the Weyl scalar is initially time independent.

1.1 What can be said about gravitational wave density value detection?

We begin with a use of particle count $n_f$ for a way to present initial GW relic inflation density using the definition given by Maggiore as a way to state that a particle count algorithm is de rigueur experimentally. And that the first place to start would be in obtaining a way to quantify $n_f$ generation in relic conditions, as a way of showing a linkage between relic GW generation, entropy, and energy flux values for the onset of inflation. We begin with

$$\Omega_{gw} = \frac{\rho_{gw}}{\rho_c} = \int_{f=0}^{f=c} d(\log f) \cdot \Omega_{gw}(f) \Rightarrow h^2_{gw} \Omega_{gw}(f) \equiv 3.6 \cdot \left( \frac{n_f}{10^{37}} \right) \left( \frac{f}{1 \text{kHz}} \right)^4$$ (1)

where $n_f$ is the frequency-based numerical count of gravitons per unit phase space.

While Maggiore’s explanation, and his treatment of gravitational wave density is very
good, the problem we have is that any relic conditions for GW involve stochastic background, and also that many theorists have fixated upon either turbulence/ and or other forms of plasma induced generation of shock waves, as stated by Duerrer\textsuperscript{2} and others looking at the electro weak transition as a GW generator. The energy flux formalism of Alcubierre\textsuperscript{3} is a natural way to obtain a criteria which we think explains how ultra low values of entropy, as given by Smoot\textsuperscript{4}, and others in the Ercole Chalonge Paris colloquia 2007, would arise in initial inflationary cosmology.

In doing so, we will expand upon a counting algorithm for entropy as given by both the author and Y. J. Ng\textsuperscript{5}, which will, when combined with an expression of energy carried per graviton complete our analysis of the relative importance of particle counting, GW energy, energy per graviton and the linkage of all these factors to initially low entropy.

The author suggests that \( n_f \) may also depend upon the interaction of gravitons with neutrinos in plasma during early-universe nucleation, as modeled by M. Marklund \textit{et al}.\textsuperscript{6}. But the main datum to consider would be in analyzing an expression given by Alcubierre’s\textsuperscript{3} formalism about energy flux, assuming that there is a solid angle for energy distribution \( \Omega \) for the energy flux to travel through.

\[
\frac{dE}{dt} = \lim_{r \to \infty} \left[ \frac{r^2}{16\pi} \right] \left| \int_{-\infty}^{t} \Psi_4 dt \right|^2 \cdot d\Omega
\]  

(2)

The expression \( \Psi_4 \) is a Weyl scalar which we will write in the form of

\[
\Psi_4 = \left[-\frac{1}{4} \left[ \partial^2_{\tau} h^+ - 2\partial_{\tau} \partial_{\tau} h^+ + \partial^2_{\tau} h^+ \right] + \frac{i}{4} \left[ \partial^2_{\tau} h^+ - 2\partial_{\tau} \partial_{\tau} h^+ + \partial^2_{\tau} h^+ \right] \right]
\]  

(4)

Our assumptions are simple, that if the energy flux expression is to be evaluated properly, before the electro weak phase transition, that time dependence of both \( h^+ \) and \( h^+ \) is miniscule and that initially \( h^+ \approx h^+ \), so as to initiate a rewrite of Eq. (4) above as

\[
\Psi_4 \approx -\frac{1}{4} \left[ \partial^2_{\tau} h^+ \right] \cdot \left( -1 + i \right)
\]  

(5)

The upshot, is that the initial energy flux about the inflationary regime would lead to looking at

\[
\left| \int_{-\infty}^{t} \Psi_4 dt \right| \approx \frac{1}{2} \left[ \partial^2_{\tau} h^+ \right] \cdot \left( \vec{n} \cdot t_{\text{Planck}} \right)
\]  

(6)
This will lead to an initial energy flux at the onset of inflation which will be presented as

$$\frac{dE}{dt} = \left[ \frac{r^2}{64\pi} \right] + \partial^2 \hat{h} + \left[ \hat{n} \cdot t_{\text{Planck}} \right]^2 \cdot \Omega \quad (7)$$

If we are talking about an initial energy flux, we then can approximate the above as

$$E_{\text{initial flux}} = \left[ \frac{r^2}{64\pi} \right] + \partial^2 \hat{h} + \left[ \hat{n} \cdot t_{\text{Planck}} \right]^2 \cdot \Omega_{\text{effective}} \quad (8)$$

Inputs into both the expression \( \partial^2 \hat{h} \), as well as \( \Omega_{\text{effective}} \) will comprise the rest of this document, plus our conclusions. The derived value of \( \Omega_{\text{effective}} \) as well as \( E_{\text{initial flux}} \) will be tied into a way to present energy per graviton, as a way of obtaining \( n_f \). The \( n_f \) value so obtained, will be used to make a relationship, using Y. J. Ng’s entropy counting algorithm of roughly \( S_{\text{entropy}} \sim n_f \). We assert that in order to obtain \( S_{\text{entropy}} \sim n_f \) from initial graviton production, as a way to quantify \( n_f \), that a small mass of the graviton can be assumed.

2. Does the graviton have small mass initially? Seeming violation of the correspondence principle. And a macro effect from small graviton mass.

We begin our inquiry by initially looking at a modification of what was presented by R. Maartens.

$$m_n(\text{Graviton}) = \frac{n}{L} + 10^{-65}\, \text{grams} \quad (9)$$

Note that Rubakov writes KK graviton representation as, after using the following normalization

$$\int \frac{dz}{a(z)} [h_n(z) \cdot h_n(z)] = \delta(m - \bar{m}) \quad \text{where } J_1, J_2, N_1, N_2 \text{ are different forms of Bessel functions, to obtain the KK graviton/ DM candidate representation along RS dS brane world}$$

$$h_n(z) = \sqrt{\frac{m}{k}} \cdot J_1(m/k) \cdot N_2\left([m/k] \cdot \exp(k \cdot z)\right) - N_1(m/k) \cdot J_2\left([m/k] \cdot \exp(k \cdot z)\right) \quad (10)$$

This Eq. (10) is for KK gravitons having a TeV magnitude mass \( M_Z \sim k \) (i.e. for mass values at .5 TeV to above a TeV in value) on a negative tension RS brane. What would
be useful would be managing to relate this KK graviton, which is moving with a speed proportional to $H^{-1}$ with regards to the negative tension brane with $h \equiv h_n(z \to 0) = const \cdot \frac{m}{\sqrt{k}}$ as an initial starting value for the KK graviton mass, before the KK graviton, as a ‘massive’ graviton moves with velocity $H^{-1}$ along the RS dS brane. If so, and if $h \equiv h_n(z \to 0) = const \cdot \frac{m}{\sqrt{k}}$ represents an initial state, then one may relate the mass of the KK graviton, moving at high speed, with the initial rest mass of the graviton, which in four space in a rest mass configuration would have a mass lower in value, i.e. of $m_{graviton}(4-Dim \ GR) \sim 10^{-48} \ eV$, as opposed to $M_X \sim M_{KK-Graviton} \sim 5 \times 10^5 \ eV$. Whatever the range of the graviton mass, it may be a way to make sense of what was presented by Dubovsky et.al. who argue for graviton mass using CMBR measurements, of $M_{KK-Graviton} \sim 10^{-20} \ eV$. Dubosky et. al. results can be conflated with Alves et al. arguing that non zero graviton mass may lead to an acceleration of our present universe, in a manner usually conflated with DE, i.e. their graviton mass would be about $m_{graviton}(4-Dim \ GR) \sim 10^{-48} \times 10^{-3} \ eV \sim 10^{65}$ grams. Also Eq. (11) will be the starting point used for a KK tower version of Eq. (12) below. So from Maarten’s 2005 paper,

$$a^2 = \left[\frac{\kappa^2}{3} \left[\rho + \frac{\rho^2}{2\lambda}\right] a^2 + \frac{\Lambda \cdot a^2}{3} + \frac{m}{a^2} - K\right]$$ (11)

Maartens also gives a 2nd Friedman equation, as

$$H^2 = -\left[\frac{\kappa^2}{2} \left[\rho + \rho^2\right] \left[1 + \frac{\rho^2}{\lambda}\right]\right] + \frac{\Lambda \cdot a^2}{3} - 2 \frac{m}{a^4} + \frac{K}{a^2}$$ (12)

Also, if we are in the regime for which $\rho \cong -P$, for red shift values $z$ between zero to 1.0-1.5 with exact equality, $\rho = -P$, for $z$ between zero to .5. The net effect will be to obtain, due to Eq. (11), and Eq. (12), and use $a \equiv a_0 = 1/(1+z)$. As given by Beckwith

$$q = -\frac{\ddot{a}}{a^2} = -1 + \frac{2}{1 + \kappa^2 \left[\rho / m\right] \cdot (1 + z)^4 \cdot (1 + \rho / 2\lambda)}$$ (13)
Eq. (11) assumes $\Lambda = 0 = K$, and the net effect is to obtain, a substitute for DE, by presenting how gravitons with a small mass done with $\Lambda \neq 0$, even if curvature $K = 0$. The ‘density’ function, $\rho$, assumed is similar to what was done by Alves et al.

2.1 Consequences of small graviton mass for reacceleration of the universe

Using Eq. (13) lead to the prediction given in Fig (1) below

![Fig. 1: Reacceleration of the universe based on Beckwith (note that $q < 0$ if $z < .423$)](image)

Now that this is presented, we should consider what the effects of a small graviton mass would be for initial entropy/ information counting at the onset of inflation.

3. Examination of Weyl scalar in the onset of inflation, to obtain entropy counting initially?

Note, from Valev, $m_{\text{graviton}}^\text{RELATIVISTIC} < 4.4 \times 10^{-22} \text{h eV / c}^2$

\[ \Lambdach_{\text{graviton}} \equiv \frac{\hbar}{m_{\text{graviton}} c} < 2.8 \times 10^{-8} \text{meters} \] (14)

This could be an argument with non zero graviton mass what to expect as a way to formulate our $n_f$ which would permit, if given a frequency range more precise ways to obtain $n_f$, so as to find a better way to use Eq. (1) more effectively. To start this, look at setting, with $\hat{A}$ a radiation constant, and using a value given by Sander's with an observationally based Friedman Equation based value of
\[ H_{\text{friedman}} = \left( \frac{T_{\text{temp}}}{c} \right)^2 \cdot \sqrt{N(T_{\text{temp}})} \cdot \sqrt{\frac{4\pi G \cdot \dot{a}}{3}} \]  

(15)

So, then that, if one uses \( N(T_{\text{temp}}) \sim 10^3 \), as opposed to an upper limit specified by Kolb, et al of, at or after electro weak, \( N(T_{\text{temp}}) \sim 10^2 \), that then

\[ h^+ \sim h^x \sim \frac{1}{2r^2} \cdot \frac{H_{\text{friedman}}}{m_{\text{Planck}}} = \frac{1}{2r^2} \cdot \left( \frac{T_{\text{temp}}}{m_{\text{Planck}}} \right)^2 \cdot \sqrt{N(T_{\text{temp}})} \cdot \sqrt{\frac{4\pi G \cdot \dot{a}}{3}} \]

(15)

For the sake of inflationary applications, we will assume that \( \sqrt{N(T_{\text{temp}})} \cdot \sqrt{\frac{4\pi G \cdot \dot{a}}{3}} \) has no spatial dependence worth speaking of, which leads to, if \( \frac{T_{\text{temp}}}{m_{\text{Planck}}} \propto \) constant of order unity. Also \( N(T_{\text{temp}}) \approx 10^2 - 10^3 \). Then one has

\[ E_{\text{initial--flux}} \sim \left[ \frac{3 \cdot N(T_{\text{temp}})}{64\pi \cdot r^6} \right] \cdot \left( \frac{\dot{n} \cdot t_{\text{Planck}}}{\Omega_{\text{effective}}} \right)^3 \]

(16)

For the sake of argument, Beckwith used \( \Omega_{\text{effective}} < 4\pi \), and experimented with values of \( r < 10^{-16} \) \( m \), and also, of \( \dot{n} \cdot t_{\text{Planck}} \approx 2 \times 10^{-35} \) sec, \( E_{\text{initial--flux}} \sim 10^{-20} \) eV. If one uses the \( m_{\text{graviton}} \sim 10^{-65} \) grams, then Beckwith obtained \( n_f \approx 10^6 \) to \( 10^7 \). This is assuming a very high initial frequency for the relic particles. We will in the next section comment upon

4. Conclusion. Examining information exchange between different universes?

Beckwith has concluded that the only way to give an advantage to higher dimensions as far as cosmology would be to look at if a fifth dimension may present a way of actual information exchange to give the following parameter input from a prior to a present universe, i.e. the fine structure constant, as given by

\[ \tilde{\alpha} = e^2 / \hbar \cdot c = \frac{e^2}{\hbar c} \times \frac{\lambda}{d} \]

(17)
The wave length as may be chosen to do such an information exchange would be part of a graviton as being part of an information counting algorithm as can be put below, namely:

Argue that when taking the log, that the $1/N$ term drops out. As used by Ng\textsuperscript{5}

$$Z_N \sim \left(1/N\right) \left(V/\lambda^3\right)^N$$

(18)

This, according to Ng\textsuperscript{5}, leads to entropy of the limiting value of, if $S = \left(\log Z_N\right)$ will be modified by having the following done, namely after his use of quantum infinite statistics, as commented upon by Beckwith\textsuperscript{7}

$$S \approx N \cdot \left(\log V/\lambda^3 \right) + 5/2 \approx N$$

(19)

Eventually, the author hopes to put on a sound foundation what ‘tHooft\textsuperscript{14} is doing with respect to ‘tHooft\textsuperscript{14} deterministic quantum mechanics and equivalence classes embedding quantum particle structures. Furthermore, making a count of gravitons with $S \approx N \sim 10^7$ gravitons\textsuperscript{7}, with Seth Lloyd’s\textsuperscript{15}

$$I = S_{\text{total}} / k_B \ln 2 = \left[\# \text{operations}\right]^{3/4} \sim 10^7$$

(20)

as implying at least one operation per unit graviton, with gravitons being one unit of information, per produced graviton\textsuperscript{7}. Note, Smoot\textsuperscript{4} gave initial values of the operations as

$$\left[\# \text{operations}\right]_{\text{initially}} \sim 10^{10}$$

(21)

The argument so presented, gives a first order approximation as to how to obtain Eq. (21)

References

1. M. Maggiore, \textit{Gravitational Waves, Volume 1: Theory and Experiment}, Oxford Univ. Press (2008)
2. R. Durrer, Massimiliano Rinaldi, \textit{Phys. Rev. D}79:063507,2009, http://arxiv.org/abs/0901.0650
3. M. Alcubierre, “Introduction to Numerical relativity”, Oxford University Press, 2008
4. G. Smoot, \textit{http://chalonge.obspm.fr/Paris07_Smoot.pdf}
5. Y. Ng, Entropy 2008, 10(4), 441-461; DOI: 10.3390/e10040441
6. M. Marklund, G. Brodin, and P. Shukla, \textit{Phys. Scr.} T82 130-132 (1999).
7. A. Beckwith, \textit{http://vixra.org/abs/0912.0012}, v 6 (newest version).
8. R. Maartens, \textit{Brane-World Gravity}, \textit{http://www.livingreviews.org/lrr-2004-7} (2004); R. Maartens \textit{Brane world cosmology}, pp 213-247 from the conference \textit{The physics of the Early Universe}, editor Papantronopoulos, (Lect. notes in phys., Vol 653, Springer Verlag, 2005).
9. V. Rubakov, Classical Theory of Gauge Fields, Princeton University press, 2002.
10. S. Dubovsky, R. Flauger, A. Starobinsky, I. Tkachev, report UTTG-06-09, TCC-23-09, \textit{http://arxiv.org/abs/0907.1658}.
11. E. Alves, O. Miranda, and J. de Araujo, arXiv: 0907.5190 (July 2009).
12. D. Valev, \textit{Aerospace Res. Bulg.} 22:68-82, 2008; \textit{http://arxiv.org/abs/hep-ph/0507255}.
13. R. Sanders, “Observational Cosmology”, pp 105-137, in “The Physics of the Early Universe,” Lecture Notes in physics, 653, E. Papantonopoulou, Editor, Springer Verlag, 2005.

14. G. ’t Hooft, http://arxiv.org/PS_cache/quant-ph/pdf/0212/0212095v1.pdf (2002); G. ’t Hooft., in Beyond the Quantum, edited by Th. M. Nieuwenhuizen et al. (World Press Scientific 2006), http://arxiv.org/PS_cache/quant-ph/pdf/0604/0604008v2.pdf, (2006).

15. S. Lloyd, Phys. Rev. Lett. 88, 237901 (2002)