CaAl$_2$Si$_2$O$_8$: Ce$^{3+}$, Tb$^{3+}$: A novel high-efficiency luminescent material for white light-emitting diodes

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Abstract. A novel rare earth activated phosphor Ca$_{1-x-y}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$, yTb$^{3+}$ were prepared under a weak reducing atmosphere by conventional solid-state reaction method and their luminescence properties under 351 nm excitation were investigated in detail. The results show that the CaAl$_2$Si$_2$O$_8$: Tb$^{3+}$ shows excited at short-UV region, the long-UV absorption peaks is very weak, which cannot be effectively excited by near ultraviolet light-emitting diodes (NUV-LEDs). However, by incorporation of Ce$^{3+}$, in the excitation spectra, strong and broad excitation bands at $\sim$ 351nm were observed, which could well match with the emission wavelength of the NUV-LEDs chip. Simultaneity, the green characteristic emission of Tb$^{3+}$ is obviously enhanced, and the emission intensity of CaAl$_2$Si$_2$O$_8$:Ce$^{3+}$, Tb$^{3+}$ phosphor is 10 times higher than that of the green phosphor CaAl$_2$Si$_2$O$_8$:Tb$^{3+}$ phosphor. The results indicates an efficient energy transfer from Ce$^{3+}$ to Tb$^{3+}$ was found and the energy transfer mechanism of Ce$^{3+}$ $\rightarrow$ Tb$^{3+}$ was studied in detail.

1. Introduction
In the global energy crisis, environments under the condition of constantly improve, the miniaturization of semiconductor light-emitting diodes (LEDs) lighting have been recognized as a new generation of lighting source due to their superior features of long life, energy saving, safety and environmental friendliness, which have attracted more and more attention of people [1, 2]. At present, the main strategy to obtain the commercialization of white light-emitting diodes (LEDs) is the combination of blue LED chip with Y$_3$Al$_5$O$_{12}$:Ce$^{3+}$ (YAG: Ce) yellow phosphor. Although this device have high luminous efficiency (the latest record has topped 300 lm/W), which is far higher than that of the traditional incandescent lamp (about 16 lm/W) and fluorescent lamp (about 70 lm/W), but this method suffers the problems of high correlated color temperature ($T_c > 5000$ K) and low color rendering index ($R_c < 70$) for general illumination due to the insufficient red-light contribution [3]. In order to solve this problem and produce white light with suitable correlated color temperature (CCT) and excellent color rendering index (CRI), trichromatic white LEDs have been proposed by pumping the tricolor (red, green and blue) phosphors with near-ultraviolet (NUV) LED chips [4]. Thus, the white light LEDs luminous efficiency and CRI strongly depends on the tricolor phosphors, the luminescence properties of phosphor plays a key role for "phosphor-converted" technology in the white light LEDs. At present, the commercial green phosphor is ZnS: Cu$^+$, Al$^{3+}$, but the sulfide phosphors has a bad stability in air, easily deliquescence, and which is easy to decompose and produce harmful gases. So the application of sulfide phosphors are
limited [5]. In recent years, nitride and nitrogen oxides have been widely studied because of their excellent luminescent properties, but the development of these phosphors is greatly restricted by the harsh preparation conditions [6]. In this regard, it is essential to develop a novel green phosphors which can be efficiently excited by the NUV-LED chips.

Over the past few decades, about the Ce$^{3+} \rightarrow$Tb$^{3+}$ energy transfer of luminescence materials have been widely research [7-9], which is mainly focused on the study of new and efficient green phosphors as the green components of low-pressure mercury vapor fluorescent lamps. Because Tb$^{3+}$ ion is a good activator with green characteristic emission, its f-f absorption in the long wave ultraviolet region is usually relatively weak and cannot be effectively excited by near ultraviolet LED chips. Considering that Ce$^{3+}$ is not only an efficient luminescent center, but also a good sensitizer, which can effectively sensitize the luminescence center of Tb$^{3+}$. In theory, there is energy transfer in Ce$^{3+} \rightarrow$Tb$^{3+}$, but the energy transfer efficiency is different in different host systems. Alkaline-earth aluminum silicate is an important mineral resource. It has been widely used in petroleum, plastics processing, building materials and light industry. In recent years, its research as the substrate of luminescent materials has attracted much attention. Alkaline-earth aluminum silicate belongs to feldspar, in which anorthite CaAl$_2$Si$_2$O$_8$ belongs to triclinic system. Essentially, the structures of alkaline-earth metal aluminum silicates are based on an infinite three-dimensional silicon–oxygen and aluminum–oxygen frameworks, which are formed from interlinked corner-sharing [SiO$_4$]$^{4-}$ and [AlO$_4$]$^{5-}$ tetrahedral [10,11]. This structure has large crystal cells and a variety of positions for different ions. Most of the research on this system is focused on the application of crystal and glass ceramics, but less on the substrate of luminescent materials. It is reported that there is a study on the luminescence properties of blue phosphor (Ca$_{1-x}$Mg$_x$)Al$_2$Si$_2$O$_8$: Eu$^{2+}$ in PDP [12], and the blue long afterglow material of CaAl$_2$Si$_2$O$_8$:Eu$^{2+}$, Dy$^{3+}$ [13, 14]. However, the energy transfer process and luminescence kinetics of Ce$^{3+} \rightarrow$Tb$^{3+}$ have not been reported.

Therefore, in this paper, CaAl$_2$Si$_2$O$_8$:Ce$^{3+}$, Tb$^{3+}$ phosphors were synthesized by high temperature solid phase method and the luminescence properties in NUV region were investigated. We focused on the concentration change and energy transfer of Ce$^{3+} \rightarrow$Tb$^{3+}$, trying to increase the absorption peak of near ultraviolet (330 ~ 410nm) through Ce$^{3+}$ to Tb$^{3+}$ energy transfer. A new and efficient green phosphor for near ultraviolet LED has been developed.

2. Experimental

2.1. Materials and synthesis

A series powder samples Ca$_{0.95-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$ (CAS: xCe$^{3+}$) (0.005 ≤ x ≤ 0.06) and Ca$_{0.95-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$, 0.05Tb$^{3+}$ (CAS: xCe$^{3+}$, 0.05Tb$^{3+}$) (0≤x≤0.06) were prepared by a traditional high temperature solid-state reaction process. Stoichiometric amounts of the starting materials CaCO$_3$ (99 %), Al$_2$O$_3$ (99 %), H$_2$SiO$_3$, CeO$_2$ (99.5 %), H$_2$BO$_3$ (99.5 %) and Tm$_2$O$_3$ (99.99 %) were weighed and blended in agate mortar thoroughly ground for 30min-1 h in alcohol. Subsequently, the dried powder mixture was placed in an alumina crucible and then preheated at 900 °C for 3 h in air, then fully ground and heated at 1300°C for 12 h in a weak reductive atmosphere of 5% H$_2$ - 95% N$_2$ gas. Finally, the prepared phosphors were cooled to room temperature and reground to get the final phosphors.

2.2. Measurements and characterization

The phase purity of all the synthesized samples were identified by a Rigaku D/MAX-2400 powder X-ray diffractometer (XRD) operating at 40 kV/60 mA with Ni-filtered Cu Kα radiation (λ = 1.54056Å). The scanning rate was seted at 10°/ min with 20 ranges from 10° to 80° with a step size of 0.02°. The Excitation (PLE) and emission (PL) spectra of the samples were measured using an FLS-920T fluorescence spectrophotometer equipped with a 450 W xenon arc lamp light source and double excitation monochromators (Edinburgh Instrument, Britain). All the measurements were performed at room temperature.
3. Results and Discussion

The XRD patterns of Ca$_{1-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$ for different Ce$^{3+}$ dopants concentration are shown in Fig. 1, along with the standard pattern of JCPDS#41-1486 as a reference. The results showed that when the doping concentration of Ce$^{3+}$ was less than 0.06, no other phase was observed, the XRD patterns of obtained samples were in good match with the CaAl$_2$Si$_2$O$_8$ phase (JCPDS No.41-1486), which has a triclinic crystal system.

Fig. 2 shows the powder XRD patterns of doubly doped series sample Ca$_{0.95-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$, 0.05Tb$^{3+}$ (0 ≤ x ≤ 0.06). It can be seen that the X-ray diffraction peak of all samples are consistent with the JCPDS standard card (No.41-1486), confirming the formation of the single-phase nature.

Fig. 3 displays the excitation and emission spectra of CaAl$_2$Si$_2$O$_8$:0.0075Ce$^{3+}$ powders. The excitation spectrum is composed of three excitation peaks under the monitoring of 419 nm. The main peaks is located at 245 nm, 284 nm and 351 nm, respectively. Among them, the strongest excitation peak is at

![Figure 1. The XRD patterns of Ca$_{1-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$ for different Ce$^{3+}$ dopants concentration](image1)

![Figure 2. XRD patterns of prepared samples and standard data of Ca$_{0.95-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$, 0.05Tb$^{3+}$ (0≤x≤0.06)](image2)

![Figure 3. The excitation (λ$_{em}$ = 419 nm) and emission spectra (λ$_{ex}$=351 nm) for CaAl$_2$Si$_2$O$_8$:0.0075Ce$^{3+}$ powders.](image3)

![Figure 4. The Emission spectra of Ca$_{1-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$. (0.005≤x≤0.06) (λ$_{ex}$=351 nm)](image4)
351 nm, which belongs to the transition from the ground state of Ce$^{3+}$ to the different splitting energy levels in the 5d excited state, and the splitting of crystal field is \(~\) 12 400 cm$^{-1}$. The emission spectrum was excited by 351 nm, shows a broad emission band with a main peak located in the blue violet region at \(~\) 419 nm (23 866 cm$^{-1}$), belonging to the 5d$\rightarrow$4f ($^2$F$_{5/2}$ and $^2$F$_{7/2}$) transition of Ce$^{3+}$. In addition, when the sample was excited at 245 nm, 284 nm or 351 nm, the shape and peak position of emission spectrum did not change, but the intensity was different. In general, the typical doublet emission bands due to the transition of the Ce$^{3+}$ ions from the 5d excited state to the $^2$F$_{5/2}$ and $^2$F$_{7/2}$ ground states, and the energy level difference is about 2000 cm$^{-1}$. However, there is only one emission band in the system, which indicates that the spin orbit coupling of Ce$^{3+}$ ground state is very weak. Besides, the emission spectrum is slightly asymmetrical towards long wave direction, indicating that the crystal field around Ce$^{3+}$ is stronger than its spin orbit coupling.

Fig.4 shows the emission spectra of Ca$_{1-x}$Al$_2$Si$_2$O$_8$:xCe$^{3+}$ ($0.005 \leq x \leq 0.06$) with different concentrations of Ce$^{3+}$ ($\lambda_{ex}=351$ nm). It is obviously observed that Ca$_{1-x}$Al$_2$Si$_2$O$_8$: xCe$^{3+}$ can be efficiently excited by UV light and emit the blue light from 5d-4f transition of Ce$^{3+}$. With the change of Ce$^{3+}$ doping concentration, the peak shape and position of emission spectra of all samples are consistent. The illustrations show the relationship between the luminescence intensity of a series of samples Ca$_{1-x}$Al$_2$Si$_2$O$_8$: xCe$^{3+}$ ($0.005 \leq x \leq 0.06$) and the doping concentration of Ce$^{3+}$. It can be seen that with the increase of Ce$^{3+}$ doping concentration, the luminescence intensity is gradually enhanced, and there is no concentration quenching phenomenon. This phenomenon can be explained that the Stokes shift of Ce$^{3+}$ is relatively large (\(~\) 4624 cm$^{-1}$) in this system, meaning that the overlap of excitation and emission spectra is very small (\(~\) 0.2 eV$^{-1}$), which indicates that the non-radiative energy transfer efficiency between adjacent luminescent center ion Ce$^{3+}$ and Ce$^{3+}$ is relatively low. The small illustrations in Fig. 4 give a photo of the sample CaAl$_2$Si$_2$O$_8$:0.06Ce$^{3+}$ with the highest luminescence intensity under the ultraviolet light, which can be seen to emit bright blue light.

![Figure 5](image1.png)

**Figure 5.** The excitation ($\lambda_{em} = 542$ nm) and emission spectra ($\lambda_{ex} = 351$ nm) for (a) CaAl$_2$Si$_2$O$_8$:0.05Tb$^{3+}$ and (b) CaAl$_2$Si$_2$O$_8$:0.06Ce$^{3+}$, 0.05Tb$^{3+}$.

![Figure 6](image2.png)

**Figure 6.** The Emission spectra of Ca$_{0.95-x}$Al$_2$Si$_2$O$_8$: xCe$^{3+}$, 0.05Tb$^{3+}$ ($0.005 \leq x \leq 0.06$) ($\lambda_{ex} = 351$ nm). The inserted figure shows the photo of CaAl$_2$Si$_2$O$_8$:0.06Ce$^{3+}$, 0.05Tb$^{3+}$ under UV.
3) of \( \text{Tb}^{3+} \), respectively, in which the \( ^5\text{D}_4 \rightarrow ^7\text{F}_5 \) transition intensity of 542nm is the largest, which is the characteristic green emission of \( \text{Tb}^{3+} \). If observed carefully, there is a very weak emission peak at 400 ~ 480 nm, which is attributed to the transition \( ^5\text{D}_3 \rightarrow ^7\text{F}_J \) \((J= 6, 5, 4, 3, 2, 1)\) of \( \text{Tb}^{3+} \). This has also been reported in other \( \text{Tb}^{3+} \) doped matrix [15-17], that is, when the doping concentration of \( \text{Tb}^{3+} \) is low, the transition emission of \( ^5\text{D}_3 \rightarrow ^7\text{F}_J \) and \( ^5\text{D}_4 \rightarrow ^7\text{F}_J \) exists simultaneously. When the doping concentration of \( \text{Tb}^{3+} \) is higher, the emission of \( ^5\text{D}_3 \rightarrow ^7\text{F}_J \) transition gradually decreases to disappear, while the main factor was \( ^5\text{D}_4 \rightarrow ^7\text{F}_J \) transition emission. This is due to the cross-relaxation between \( ^5\text{D}_3 \) and \( ^5\text{D}_4 \) [18], when the \( \text{Tb}^{3+} \) doping concentration is high, the average distance between \( \text{Tb}^{3+} \) and \( \text{Tb}^{3+} \) becomes shorter and shorter, and the energy will transfer from the \( ^5\text{D}_3 \) level of one \( \text{Tb}^{3+} \) to the \( ^5\text{D}_4 \) level of another \( \text{Tb}^{3+} \). Thus, the \( ^5\text{D}_3 \rightarrow ^7\text{F}_J \) transition emission is quenched, and the cross-relaxation process can be expressed as follows:

\[
\text{Tb}^{3+}(^5\text{D}_3)+ \text{Tb}^{3+}(^7\text{F}_0) \rightarrow \text{Tb}^{3+}(^5\text{D}_4)+ \text{Tb}^{3+}(^7\text{F}_0)
\]

From the excitation and emission spectra of \( \text{Ce}^{3+} \), \( \text{Tb}^{3+} \) double doped samples (b), it is shown that in the excitation spectra that under the monitoring of \( ^5\text{D}_3 \rightarrow ^7\text{F}_J \) transition emission (542 nm) of \( \text{Tb}^{3+} \), there are not only the excitation peaks of \( \text{Ce}^{3+} \) but also the excitation peaks of \( \text{Tb}^{3+} \). This indicates that \( \text{Tb}^{3+} \) is essentially excited by \( \text{Ce}^{3+} \), implying that \( \text{Ce}^{3+} \rightarrow \text{Tb}^{3+} \) has an effective energy transfer, and that the excitation peak of \( \text{Ce}^{3+} \) at ~ 351 nm is the strongest, which matches the emission spectrum of NUV-LED chip. The emission spectrum excited at 351 nm consists of two parts: the broadband emission of the 5d\( \rightarrow \)4f transition of \( \text{Ce}^{3+} \) (400 ~ 480 nm) and the peak emission of the \( ^5\text{D}_4 \rightarrow ^7\text{F}_J \) \((J= 6, 5, 4, 3)\) of \( \text{Tb}^{3+} \). From the characteristic transition emission intensity of \( \text{Tb}^{3+} \) in the two diagrams (a) and (b), it can be seen that the emission intensity of co-doped \( \text{Tb}^{3+} \) is 10 times that of single-doped \( \text{Tb}^{3+} \). It is obvious that the significant enhancement of \( \text{Tb}^{3+} \) emission intensity is the result of \( \text{Ce}^{3+} \rightarrow \text{Tb}^{3+} \) energy transfer.

Fig. 6 shows the emission spectra of Ca\(_{0.95-x}\)Al\( _2\)Si\( _2\)O\( _8\): \( x\text{Ce}^{3+}, 0.05\text{Tb}^{3+} \) \((0 \leq x \leq 0.06)\) excited at 351 nm. The fixed concentration of \( \text{Tb}^{3+} \) in all samples was 0.05 mol, and the doping concentration of \( \text{Ce}^{3+} \) was changed. It can be seen from the diagram that with the increase of \( \text{Ce}^{3+} \) doping concentration, the emission intensity of \( \text{Ce}^{3+} \) gradually weakened while the emission intensity of \( \text{Tb}^{3+} \) increased gradually, which indicates that the introduction of \( \text{Ce}^{3+} \) has a process of energy transfer to \( \text{Tb}^{3+} \). In this process, \( \text{Tb}^{3+} \) obtains partial excitation energy from \( \text{Ce}^{3+} \), resulting in the enhancement of the \( ^5\text{D}_4 \rightarrow ^7\text{F}_J \) \((J= 6, 5, 4, 3)\) transition emission of \( \text{Tb}^{3+} \). In addition, with the increase of \( \text{Ce}^{3+} \) concentration, the distance between \( \text{Ce}^{3+} \) and \( \text{Tb}^{3+} \) is getting closer and closer, and the energy transfer between the two ions increases gradually. Therefore, the emission intensity of \( \text{Tb}^{3+} \) is gradually enhanced. The small illustrations in Fig. 6 shows the photo of the most luminous intensity sample CaAl\( _2\)Si\( _2\)O\( _8\): \( 0.06\text{Ce}^{3+}, 0.05\text{Tb}^{3+} \) irradiated by ultraviolet lamp. It can be seen that the sample emits very bright green light.

The above is a qualitative analysis of the possible energy transfer process in the system based on the experimental data. The mechanism of energy transfer will be studied below. The energy transfer between the sensitizer and the activator usually has three forms: radiation transfer (radiative reabsorption), electron cloud exchange, and the effect of the near field force of the multipole. The radiation transfer process is that if the energy of the radiation spectrum emitted by one ion coincides with that of another ion absorption spectrum, then the radiation light will be absorbed by another ion, and the energy transfer process will occur between the ions. That is, the transfer process of "radiation reabsorption". For the process of radiation reabsorption, the emission spectrum of sensitizer and the excitation spectrum of activator are overlapped or partially overlapped, but for rare earth ions, the f-f transition is linear, both emission and absorption intensity are relatively weak. Therefore, the energy transfer efficiency of the radiation reabsorption process is relatively low. In this system, there is a small overlap between the emission peaks of \( \text{Ce}^{3+} \) and the f-f transition absorption line of \( \text{Tb}^{3+} \), so the radiation energy transfer between the sensitizer and the activator can be ignored. In the case of electron cloud exchange, there is a very close distance between the activator and the sensitizer (< 4 Å), which means that there is a large overlap between the two electron cloud orbits. However, both \( \text{Ce}^{3+} \) and \( \text{Tb}^{3+} \) are reduced ions in this system. If the electron cloud exchange is to take place, a higher energy is needed [19]. It can be seen that it is impossible to achieve energy transfer by the action of the electron cloud exchange. Therefore,
the energy transfer mechanism of Ce\(^{3+} \rightarrow \)Tb\(^{3+}\) in CaAl\(_2\)Si\(_2\)O\(_8\):Ce\(^{3+}\), Tb\(^{3+}\) system is mainly due to the action of the near field force of non-radiative multipole.

4. Conclusion
In the CaAl\(_2\)Si\(_2\)O\(_8\) system, the luminescent properties of single doped Ce\(^{3+}\) and Tb\(^{3+}\) samples and the energy transfer mechanism of Ce\(^{3+} \rightarrow\)Tb\(^{3+}\) in the double doped samples were studied. For the sample doped with Ce\(^{3+}\), it has a strong absorption in the near ultraviolet region of 351nm. Its emission spectrum is a broad band with a main peak at 419 nm, which belongs to the 5d\(^{0} \rightarrow \)4f of Ce\(^{3+}\). It emits very strong blue light. The highest luminescence intensity of the sample is CaAl\(_2\)Si\(_2\)O\(_8\):0.06Ce\(^{3+}\). However, the absorption of Tb\(^{3+}\) doped samples in the near ultraviolet region is very weak and cannot be effectively excited by NUV-LED. Through Ce\(^{3+}\), Tb\(^{3+}\) double doped samples, the green emission of Tb\(^{3+}\) is obviously enhanced, which is 10 times the luminescence intensity of the single doped Tb\(^{3+}\) sample, indicating that there is an effective energy transfer between Ce\(^{3+}\) and Tb\(^{3+}\), and the mechanism of Ce\(^{3+} \rightarrow\)Tb\(^{3+}\) energy transfer is studied, and the results show that there is an effective energy transfer between Ce\(^{3+}\) and Tb\(^{3+}\). Ce\(^{3+} \rightarrow\)Tb\(^{3+}\) energy transfer mechanism is mainly due to the non-radiative multipole near-field force.

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