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Single-Molecule Magnets

A Redox-Active Bridging Ligand to Promote Spin Delocalization, High-Spin Complexes, and Magnetic Multi-Switchability

Xiaozhou Ma, Elizaveta A. Suturina, Siddhartha De, Philippe Négrier, Mathieu Rouzières, Rodolphe Clérac,* and Pierre Dechambenoit*

Abstract: A dinuclear $\text{Co}^{II}$ complex, $[\text{Co}_2(tphz)(tpy)_2]^{3+}$ ($n = 4, 3$ or $2$; tphz: tetrapyridophenazine; tpy: terpyridine), has been assembled using the redox active and strongly complexing tphz bridging ligand. The magnetic properties of this complex can be tuned from spin crossover with $T_{1/2} \approx 470 \text{ K}$ for the pristine compound ($n = 4$) to single molecule magnet with an $S_T = 5/2$ spin ground state when once reduced ($n = 3$) to finally a diamagnetic species when twice reduced ($n = 2$). The two successive and reversible reductions are concomitant with an increase of the spin delocalization within the complex, promoting remarkably large magnetic exchange couplings and high spin species even at room temperature.

One of the most attractive features of molecule based materials is the ability to fine tune and control their physical properties by a suitable choice of the molecular components.[1] In magnetism, this ability is of great interest as the magnetic properties usually result from cooperative effects between the spin carriers. The choice of the linker between spin carriers is therefore crucial as it mediates the communication/interactions between them. Using redox active bridging ligands is thus particularly appealing as they are able to favor electronic conductivity, provide multistability, and allow a certain tuning of the magnetic properties by simple oxidation or reduction.[2] In particular, the isolation of bridging ligands in their stable radical forms can be seen as the key to design new materials with strong magnetic exchange couplings, as the radical linker acts as a very efficient magnetic relay between the metallic spin carriers.[3] This strategy was successfully used to synthesize single molecule magnets (SMMs), single chain magnets with high blocking temperatures,[3a–c,4] as well as the first room temperature molecule based magnet.[5]

In our quest to design an ideal redox active bridging ligand, tetrapyridophenazine (tphz) has been chosen (Supporting Information, Scheme S1). This ligand has been reported to be redox active when coordinated to Ru metal ions[6] with two successive well separated (by 0.6 V) one electron reduction processes. The reduced forms were not further investigated and, to the best of our knowledge, this ligand has never been involved in magnetically interesting materials. The strong coordination ability (thanks to a chelating bis(tridentate) coordination mode), the planarity and the extended π system of tphz are also expected to support a large delocalization of the unpaired electron when reduced once, providing an excellent radical linker candidate for the design of strongly coupled magnetic architectures.[6] To probe this promising idea, prototype dinuclear complexes have been synthesized by reaction of the bridging tphz ligand with paramagnetic metal centers capped by terpyridine (tpy) blocking ligands. Herein, we report a novel redox active dinuclear $\text{Co}^{II}$ system, $[\text{Co}_2(tphz)(tpy)_2](\text{PF}_6)_2$ ($x = 4, 3$, or 2), displaying an increasing spin delocalization upon two successive reduction processes, resulting in the promotion of remarkably large magnetic couplings, such that the $S_T = 5/2$ and $S_T = 0$ ground states of the mono and twice reduced species remain well isolated from excited states even at room temperature.

$[\text{Co}_2(tphz)(tpy)_2](\text{PF}_6)_2\cdot3\text{CH}_3\text{CN}$ (1) was prepared by reacting $\text{Co}^{II}\text{Cl}_2$[7] with tphz in the presence of TIPF$_6$ in acetonitrile. Subsequent diffusion of diethyl ether vapor into this dark red solution yields needle shaped orange crystals of 1 (see the Supporting Information). This compound was characterized by single crystal X ray diffraction at 120 K (Figure 1; Supporting Information, Table S1), confirming the envisioned structure with two crystallographically equivalent Co(tpy)$^{2+}$ units connected by tphz. The complex 1 crystallizes in the $C2/c$ space group with a two fold axis passing through its center perpendicularly to the mean plane of the tphz ligand. Each Co ions reside in an octahedral environment involving three nitrogen atoms from the bridging ligand and another three nitrogen atoms from capping terpyridine. The Co N bond distances between the metal and tpy (Co N4,N5,N6) range from 1.87 to 1.99 Å, indicating a Co$^{II}$ center in its low spin (LS) state.[9] The Co N$_{py}$ bond distance (Co N2) between Co and tphz is 1.94 Å while the other two Co N$_{py}$ distances (Co N1 and Co N3) are significantly longer due to the rigidity of the bridging ligand (Table 1).
Table 1: Selected bond distances [Å] in 1-4 complexes.

|        | 1   | 2   | 3   | 4   |
|--------|-----|-----|-----|-----|
| Co N1  | 2.353(6) | 2.344(6) | 2.396(4) | 2.650(4) |
| Co N2  | 1.935(5) | 1.970(5) | 1.923(3) | 1.976(5) |
| Co N3  | 2.286(6) | 2.320(6) | 2.269(4) | 2.144(4) |
| Co N4  | 1.991(6) | 2.129(6) | 2.152(3) | 2.001(4) |
| Co N5  | 1.875(5) | 2.089(5) | 2.069(3) | 1.885(4) |
| Co N6  | 1.947(6) | 2.148(6) | 2.128(3) | 2.008(4) |
| C C (pyz) | 1.405(8) | 1.397(8) | 1.381(6) | 1.378(8) |
| C N (pyz) | 1.331(8) | 1.356(8) | 1.375(6) | 1.392(7) |
| N2–N2  | 2.763(1) | 2.821(7) | 2.856(4) | 2.892(6) |

[a] pyz pyrazine part of tphz. [b] Average distance.

To investigate the redox properties of this prototype dinuclear complex, cyclic voltammetry of 1 was performed in acetonitrile, revealing four redox processes, centered at 0.34 V, 1.04 V, 1.77 V and 1.93 V vs. FeCl₃/FeCl₂ (Figure 2; Supporting Information, Figure S3). The rest potential is close to 0 V, indicating the presence of only reduction processes, each of them involving one electron. [5] These potentials are very close to what has been reported for the Ru analogue, [6] suggesting that the redox processes are mainly centered on the ligands, which will be further confirmed by crystal structure analysis and DFT calculations (see below). The first two signals correspond to the reduction of the central tphz whereas those at about 2 V correspond to those of the tpy ligands. The dissymmetry of the cathodic signal is very diffusion dependent and was not further investigated. The potentials of the three first reduction processes are well separated (by 0.7 V), indicating that the mono- and doubly reduced species are particularly stable. It is also remarkable to observe the presence of two peaks for the tpy ligands, indicating a significant electronic coupling through the [Co(tphz)Co] moiety.

Based on these redox properties, 1 was reacted with a slight excess of one, two, and four equivalents of a strong reductant, KC₈, affording after recrystallization the once, twice, and four times reduced analogues as [Co₂(tphz)(tpy)](PF₆)₃·2CH₃CN (2), [Co₂(tphz)(tpy)](PF₆)₂·2CH₃CN (3), and [Co₂(tphz)(tpy)]·2THF (4). Owing to the extreme instability of 4, only its single crystal X-ray analysis is reported in the following.

The overall molecular structure of the dinuclear complexes in 2, 3, and 4 is similar to 1 (Supporting Information, Figure S1, Table S1); that is, each Co center is surrounded by six nitrogen atoms from tpy and tphz ligands, each Co being equivalent by symmetry as a consequence of an inversion center or two fold axes. Examination of the bond length reveals however noticeable differences and features (Table 1; Supporting Information, Table S2). The average Co-N distances between Co and tpy ligands are 2.12 Å in 2 and 3, consistent with a high spin (HS) Co³⁺ center. In 4, this average bond distance is 1.96 Å and a clear Jahn Teller distortion is observed along the N1-Co-N3 axis, similar to 1 and what is expected for a low spin (LS) state. [9] In these complexes, the spin state assignment of the Co sites is also confirmed by DFT (B3LYP/def2-TZVP) calculations, which indicate a LS state significantly lower in energy than the HS state in 1 and 4, and contrariwise in 2 and 3 (Supporting Information, Figure S15). This is further supported by ligand field analysis and by the complete active space self consistent field method (CASSCF) carried out in ORCA 4.0, [10] which shows that the cobalt centers in 1 and 4 have larger ligand field splitting than in 2 and 3, explaining therefore their LS versus HS preferences (Supporting Information, Figure S16). The possibility to be in presence of Co⁵⁺ in 1 and Co⁷⁺ in 3 or 4 was also considered, but excluded after combined analysis of the magnetic susceptibility, crystal structure, and DFT data. Upon reduction, the evolution of C N and C C bond distances in tpy and tphz ligands (Table 1; Supporting Information, Table S2) is also confirming the previous assignment of the reduction processes. Whereas the two first reductions do not affect the bond distances in the tpy capping ligands, a significant elongation of their C N bond distances occurs in 4, sign of their reduction in the two most cathodic reduction processes (Supporting Information, Table S2). One can also notice that the two first processes are centered on the pyrazine part of tphz bridging ligand with a significant elongation of pyrazine C N bonds and a large increase of N2–N2 distance, whereas the bond distances of the pyridine moieties in tphz and tpy remain almost unchanged (Table 1). While the Co N₃ distances (Co N1,N3) remain rather long owing to the rigidity of tphz
in this series of complexes, another interesting structural feature is the remarkably short Co-N \textsubscript{pz} (Co N2) bond distance in 2 (1.97 Å) and even more in 3 (1.92 Å; Table 1). These bond distances are much shorter than what has been reported for HS Co\textsuperscript{II} centers with similar tetrakis(2 pyridyl) pyrazine (tppz) ligands\textsuperscript{[13]} In the present case, the persistently short Co N2 bond length upon reduction is likely a consequence of the enhanced electrostatic interactions upon successive reductions, which compensate the expected bond lengthening owing to the spin state switching. Those short bond distances also suggest a particularly good overlap between the Co and tphz orbitals, which could enable stronger magnetic exchange coupling than in the tpzp based systems, which are already known to mediate significantly stronger magnetic coupling.\textsuperscript{[14,11,12]}

To evaluate and compare the strength of the magnetic coupling in this family of complexes, de magnetic susceptibility measurements were performed for 1, 2, and 3 under a 0.1 T applied field. For 1, the $\chi T$ product reaches a plateau at about 100 200 K with a value of 0.85 cm\textsuperscript{3} K mol\textsuperscript{-1}, in agreement with two isolated LS Co\textsuperscript{II} centers with $g=2.13(5)$ (Figure 3). Upon cooling below 100 K, the $\chi T$ value decreases down to 0.04 cm\textsuperscript{3} K mol\textsuperscript{-1} at 1.85 K, which can be ascribed to the presence of an antiferromagnetic coupling between the two $S=1/2$ Co\textsuperscript{II} magnetic centers. Although weak, this exchange constant of 12.7(3) K (green line in Figure 3; with $g=2.17(5)$)\textsuperscript{[13]} is remarkably larger than what is usually observed through diamagnetic pyrazine bridging ligand.\textsuperscript{[13,14]}

Above 200 K, the compound undergoes a spin crossover (SCO) behavior. The $\chi T$ value at 400 K (1.72 cm\textsuperscript{3} K mol\textsuperscript{-1}) highlights an incomplete $S=1/2$ (LS) to 3/2 (HS) thermal conversion for each cobalt ion\textsuperscript{[14]} which was fitted to an ideal solution model (red line in Figure 3) assuming a constant $g$ value of 2.18 and giving rise to $T_{1/2} \approx 467(5)$ K with an associated enthalpy variation of 27(2) kJ mol\textsuperscript{-1}.\textsuperscript{[13,14]} Upon further cooling, the magnetic properties are nearly reversible but the partial loss of interstitial solvent molecules around 400 K prevents obtaining the exact same SCO behavior (Supporting Information, Figure S8).

For the monoreduced compound 2, the $\chi T$ product found at room temperature is 6.3 cm\textsuperscript{3} K mol\textsuperscript{-1}, which is much larger than the expected spin only value for two magnetically isolated $S=3/2$ Co\textsuperscript{II} and one $S=1/2$ radical (Figure 3). This result is however more consistent with an $S_{T}=1/2$ ground state (with $g=2.40(5)$) expected in presence of an extremely large exchange coupling between Co\textsuperscript{II} magnetic centers and the bridging tphz$^-$ radical spin. The $\chi T$ values remain almost constant in the 100 300 K temperature range also indicating that the $S_{T}=1/2$ ground state is well isolated from excited states, as confirmed by DFT calculations that situate the first excited quartet state at about 500 K above (Supporting Information, Figure S18). The large magnitude of the $J_{Co\text{ tphz}}$, interactions and the resulting temperature independence of the $\chi T$ product above 100 K preclude any exact quantification of $J_{Co\text{ tphz}}$ \textsuperscript{,} nevertheless it was roughly estimated to be at least 500 K by simulating the experimental susceptibility data using a three spins model with PHIL\textsuperscript{[17,18]}.

This approximated value constitutes unequivocally one of the highest interaction reported to date\textsuperscript{[3]} and is in good agreement with DFT results (509 K). Note that below 300 K, the compound is better described as a single $S_{T}=1/2$ macro spin, with a spin density delocalized on the Co centers, its coordination sphere and the whole pyrazine part of the tphz ligand (Figure 4).

The magnetic anisotropy, revealed by the decrease of the $\chi T$ product below 40 K, was estimated to be $D/k_B \approx 27.6(9)$ K \textsuperscript{[18,19]} in very good agreement with the NEVPT2 results ($D/k_B = +27$ K).

When further reducing the complex, magnetic susceptibility measurements of the resulting compound 3 reveal a continuous increasing of the $\chi T$ product from 0.07 cm\textsuperscript{3} K mol\textsuperscript{-1} at 1.8 K to 2.01 cm\textsuperscript{3} K mol\textsuperscript{-1} at 400 K without

**Figure 3.** Temperature dependence of the $\chi T$ product for 1, 2, and 3 at 0.1 T (black, purple, and blue circles, respectively). The lines represent the best fits as discussed in the text.\textsuperscript{[13,11,12]} Inset: magnetization versus field curves for 2 below 8 K.

**Figure 4.** View of the computed spin density of the broken symmetry solution for the complexes in 1, 2, and 3. Colors represent different signs of the spin density.
reaching a plateau in this temperature range (Figure 3). This thermal behavior is consistent with a strong antiferromagnetic coupling between the two HS Co(I) magnetic centers leading to a singlet ground state. Indeed, the best fit of the $\chi T$ vs. $T$ data gives rise to $J/k_B = 74(5)$ K (with $g = 2.03(5)$; blue line in Figure 3), a much larger coupling than what one could initially expect through a formal diamagnetic ligand. As a consequence of this large coupling, the singlet ground state is quite well isolated from the excited states, as clearly supported by DFT calculations ($J/k_B = 137$ K, see below). This result is also observed in solution, as confirmed by $^1$H NMR at room temperature, which shows only slightly paramagnetically shifted resonance peaks of tpy and tphz (between 7.4, 8.7 and 8.3, 11.7 ppm, respectively; Supporting Information, Figure S6), in marked contrast to 1 and 2 (Supporting Information, Figures S4, S5). The estimation of the magnetic susceptibility in solution by Evan’s method gives rise to a comparable $\chi T$ value at 300 K (2.48(8) cm$^3$K mol$^{-1}$) than in solid state (1.77(8) cm$^3$K mol$^{-1}$; see details in the Supporting Information).

Ac susceptibility measurements were also performed to probe the magnetization dynamics of these complexes and in particular the possible single molecule magnet (SMM) properties.$^{[20]}$ Frequency and temperature dependent ac signals were revealed under a small applied dc field (600 Oe; Supporting Information, Figures S9, S12) only for 2. To probe the effect of the applied dc field on the slow relaxation of magnetization, the characteristic time, $r$, was extracted by fitting both variable field and temperature ac susceptibility data with the generalized Debye model. The best fit of $r(T, H)$ was obtained considering only Direct and Raman relaxation mechanisms (Supporting Information, Figure S13), which are thus at the origin of the SMM properties of 2.$^{[21,22]}

To explain the large intramolecular magnetic interaction and the resulting well isolated ground state of $[\text{Co}_2(\text{tphz})(\text{tpy})]^{2+}$ and $[\text{Co}_2(\text{tphz})(\text{tpy})]^{3+}$ in 2 and 3 respectively, comprehensive DFT studies were performed (see details in the Supporting Information). For both complexes, the spin density calculated with the broken symmetry solution shows a significant spin delocalization on the $\pi$ system of the tphz bridging ligand as depicted on Figure 4. Indeed, as shown by ab initio ligand field analysis (Supporting Information, Figure S16), unpaired electrons of the HS Co(I) are localized on $d_{\pi,-\pi}$, $d_\pi$, and $d_\sigma$ orbitals. Whereas the $d_{\pi,-\pi}$ and $d_\pi$ orbitals overlap with nitrogen $\sigma$ orbitals, the $d_\sigma$ orbitals interact strongly with $\pi$ orbitals of tphz bridging ligand, providing a privileged channel for significant spin delocalization (Figure 4; Supporting Information, Figures S19 and S20, Table S4), and thus a strong magnetic exchange. This spin delocalization through the $\pi$ system is even slightly larger in 3 than in 2 (14 vs. 10%; Supporting Information, Table S4). These orbital considerations are often overlooked for the design of molecule based magnets except in the case of the cyanide bridge for which this approach has been extensively discussed.$^{[23]}$ As a general synthetic strategy, the present results illustrate the importance of using metal ions with singly occupied $d_\sigma$ orbital, to design highly magnetically coupled systems with aromatic bridging ligands. The large orbital delocalization is also supported experimentally by a broad absorption band in the NIR region at 782 nm for 2, and 1250 nm for 3 (Supporting Information, Figure S7), indicative of an unusually large electronic delocalization for non mixed valence systems. This ability of the bridging ligand to mediate a strong communication between the metal centers results therefore from a remarkably strong metal ligand coordination (as illustrated by the unusually short Co N2 bond) intrinsically linked to a strong overlap of $d$ and $\pi$ orbitals and an extended delocalization through the $\pi$ system.

In summary, a new multi switchable complex has been rationally assembled using a redox active and strongly complexing bridging ligand. By simple oxidation/reduction, its magnetic properties can be tuned from a spin crossover complex, to a single molecule magnet with an $S_\text{t} = 5/2$ spin ground state (when once reduced) and to diamagnetic species when twice reduced. This work illustrates experimentally and theoretically how successive redox processes can increase the spin delocalization, and thus promote a dramatic enhancement of the intramolecular magnetic coupling through a bridging ligand, not only in its radical form. This general synthetic approach should help for the intentional design of new high spin complexes and SMMs, but it could also be extended to systems with higher nuclearities and dimensionalties, for example to obtain high temperature molecule based magnets.

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Conflict of interest

The authors declare no conflict of interest.

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Cyclic voltammograms of
potential is different, 0.1 V
versus 0.8 V and 1.1 V respectively,
confirming the one electron reduction processes (Supporting Information, Figure S3).

The thermal dependence of the spin crossover was fitted to the following equation deduced from the ideal solution model:

\[ \chi^2 T = \chi T_{1S}^2 + (q_1 T_{1S})^2 + (q_2 T_{1S})^2 \]

with \( \chi T_{1S} \) and \( \chi T_{2S} \) being the \( \chi T \) products for pure LS and HS species assuming a \( g \) value of 2.18, \( B \) the ideal gas constant and \( \Delta H \) the associated enthalpy variation.

The relation \( r^1 = r_{\text{mean}} + r_{\text{isol}} \) was used to fit simultaneously the \( r \) vs. \( T \) and \( r \) vs. \( H \) data (see details in the Supporting Information).

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Jk0 and \( g \) were estimated by fitting the experimental data to the susceptibility calculated from the van Vleck equation in the weak field approximation relative to the following isotropic spin Hamiltonian:

\[ \mathbf{H} = \frac{1}{2} J \mathbf{S} \cdot \mathbf{S} \]