A NOTE ON THE LOGARITHMIC \((p,p')\) FUSION

A.M. SEMIKHATOV

ABSTRACT. The procedure in [Fuchs et al.] to obtain a fusion algebra from the modular transformation of characters in logarithmic conformal field models is extended to the \((p,p')\) logarithmic models. The resulting fusion algebra coincides with the Grothendieck ring of the quantum group of the \((p,p')\) model.

1. Introduction

This paper is a remark on fusion in a class of logarithmic models of conformal field theory [1, 2, 3]. In rational conformal field models, fusion is related to modular transformations of characters by the celebrated Verlinde formula [4, 5]. Because the Verlinde formula relies on the fact that the fusion algebra is semisimple, it does not immediately extend to logarithmic conformal field theories, where fusion algebras (starting with the pioneering results in [6]) are typically nonsemisimple. The known extensions of the Verlinde formula to the nonsemisimple realm rely on some extra input, in one form or another [7] (also see [8]). In the prescription proposed in [9], this extra input can be related to a quantum-group formulation.

The role of quantum groups in logarithmic conformal field theory gradually emerged in [10, 11, 12, 13] (see [14] for a summary and [15] for some further development), leading to a version of the Kazhdan–Lusztig “duality” between the extended algebra \(W\) in a logarithmic conformal field model and the corresponding quantum group \(g\).\(^1\) The most remarkable result related to the Kazhdan–Lusztig duality is the coincidence of modular group representations (the one generated from the \(W\) characters and the one carried by the center of \((g\)) also, the Grothendieck ring of \((g\) is a natural candidate for the fusion algebra of \(W\)-representations (we speak of the \(K_0\)-type fusion, see [9, 18]).

For the \((p,1)\) logarithmic models, in particular, this “quantum-group candidate fusion” coincides with the fusion derived in [9] from the characters, thus lending additional support to the procedure proposed in [9]. The aim of this paper is to extend the existing state of consistency to \((p,p')\) logarithmic models: we propose a prescription whereby the modular transformations of the characters of the extended algebra in the \((p,p')\) logarithmic

\(^1\)These are factorizable ribbon quantum groups at even roots of unity; see [16] for their other use and [17] for an interesting precursor of their occurrence in logarithmic models: the ribbon structure, the (co)integral, and the \(M\)onodromy matrix (cf. [10, 13]) are already present in [17], albeit in a somewhat simpler situation.
model [12] are converted into a nonsemisimple fusion algebra, which turns out to coincide with the Grothendieck ring of the corresponding quantum group \( g \) [13]. For this, we follow the approach in [9] (also see [7]) very closely. In Sec. 2, we describe our starting point, the modular group representation generated from the characters of the extended algebra of the \((p, p')\) logarithmic models. In Sec. 3, we formulate the procedure to convert these modular transformations to the following fusion algebra on \( 2pp' \) elements \( \mathcal{K}^\pm_{r,r'} \) [13]:

\[
\mathcal{K}^\alpha_{r,r'} \mathcal{K}^\beta_{s,s'} = \sum_{u=|r-s|+1}^{r+s-1} \sum_{u'=|r'-s'|+1}^{r'+s'-1} \mathcal{K}^\alpha\beta_{u,u'},
\]

where \( \alpha, \beta = \pm 1 \) and

\[
\mathcal{K}^\alpha_{r,r'} = \begin{cases} 
\mathcal{K}^\alpha_{r,r'}, & 1 \leq r \leq p, \quad 1 \leq r' \leq p', \\
\mathcal{K}^\alpha_{2p-r,r'} + 2\mathcal{K}^-\alpha_{r-p,r'}, & p+1 \leq r \leq 2p-1, \quad 1 \leq r' \leq p', \\
\mathcal{K}^\alpha_{r,2p'-r'} + 2\mathcal{K}^-\alpha_{r,rp'-1}, & 1 \leq r \leq p, \quad p' + 1 \leq r' \leq 2p'-1, \\
\mathcal{K}^\alpha_{2p-r',2p'-r'} + 2\mathcal{K}^-\alpha_{p-r',r'}, & 1 \leq r \leq 2p-1, \quad 1 \leq r' \leq 2p'-1, \\
+ 2\mathcal{K}^-\alpha_{r-p,2p'-r'} + 4\mathcal{K}^-\alpha_{r-p,r'-p'}, & p+1 \leq r \leq 2p-1, \quad p' + 1 \leq r' \leq 2p'-1.
\end{cases}
\]

The identity of this associative commutative algebra is given by \( \mathcal{K}^+_{1,1} \). We also recall from [13] that this algebra is generated by two elements \( \mathcal{K}^+_{1,2} \) and \( \mathcal{K}^+_{2,1} \) and can also be described as the quotient of \( \mathbb{C}[x,y] \) by the ideal generated by the polynomials

\[
U_{2p+1}(x) - U_{2p-1}(x) - 2, \\
U_{2p'+1}(y) - U_{2p'-1}(y) - 2, \\
U_{p+1}(x) - U_{p-1}(x) - U_{p'+1}(y) + U_{p'-1}(y),
\]

where

\[
U_s(2\cos t) = \frac{\sin st}{\sin t}, \quad s \geq 1,
\]

are Chebyshev polynomials of the second kind.

2. MODULAR TRANSFORMATIONS OF THE \((p, p')\) CHARACTERS [12]

For each pair of coprime positive integers \( p, p' \), the extended algebra of the logarithmic \((p, p')\) model is the \( W \)-algebra \( \mathcal{W}_{p, p'} \) identified and studied in [12]. It has \( 2((p-1)(p' - 1) + 2pp' \) irreducible representations, the \( \frac{1}{2}(p-1)(p' - 1) \) of which are just the Virasoro representations in the corresponding \((p, p')\) minimal model and the other are “genuine” \( \mathcal{W}_{p, p'} \)-representations (such that the radical of \( \mathcal{W}_{p, p'} \) acts nontrivially). In what follows,
the characters of irreducible $W_{p,p'}$-representations are denoted as

\[ \chi_{r',r}(\tau), \quad \chi_{r',r}^+(\tau), \quad \chi_{r',r}^-(\tau), \quad (r,r') \in \mathcal{J}_0, \quad 1 \leq r \leq p, \quad 1 \leq r' \leq p' \]

where we introduce the index set

\[ \mathcal{J}_0 = \{(r,r') \mid 1 \leq r \leq p - 1, \quad 1 \leq r' \leq p' - 1, \quad p'r + pr' \leq pp'\}, \]

with $|\mathcal{J}_0| = \frac{1}{2}(p-1)(p'-1)$ (we recall the well-known symmetry $\chi_{r',r}(\tau) = \chi_{p-r,p'-r}(\tau)$ of the minimal-model Virasoro characters).

The modular (specifically, $S$-) transformation properties of the characters are as follows. First, the minimal-model characters $\chi_{r',r}$ are well-known to $S$-transform as

\[ \chi_{r',r'}(-\frac{1}{\tau}) = -\frac{2\sqrt{2}}{\sqrt{pp'}} \sum_{(s,s') \in \mathcal{J}_0} (-1)^{rs'+sr'} \sin \frac{\pi pr's'}{p} \sin \frac{\pi pr's'}{p'} \chi_{s,s'}(\tau), \quad (r,r') \in \mathcal{J}_0. \]

Next, it follows from [12] that (for $1 \leq r \leq p$ and $1 \leq r' \leq p'$)

\[ \chi_{r',r}^+(\tau) = \sum_{s=1}^{p} \sum_{s'=1}^{p'} S_{r',s,s'}^+(\tau) \left( \chi_{s,s'}^+(\tau) + (-1)^{p'p+pr'} \chi_{s,s'}^{-}(\tau) \right) + \sum_{(s,s') \in \mathcal{J}_0} \tilde{S}_{r',s,s'}^+(\tau) \chi_{s,s'}(\tau), \]

\[ \chi_{r',r}^-(\tau) = \sum_{s=1}^{p} \sum_{s'=1}^{p'} (-1)^{ps'+ps'} S_{r',s,s'}^{-}(\tau) \left( \chi_{s,s'}^+(\tau) + (-1)^{p'p+pr'} \chi_{s,s'}^{-}(\tau) \right) + \sum_{(s,s') \in \mathcal{J}_0} \tilde{S}_{r',s,s'}^{-}(\tau) \chi_{s,s'}(\tau), \]

where the matrix elements $S_{r',s,s'}(\tau)$ that interest us in what follows are given by

\[ S_{r',s,s'}(\tau) = \frac{2\sqrt{2}}{\sqrt{pp'}} (-1)^{rs'+sr'} \left( \frac{r}{p} \cos \frac{\pi pr's'}{p} - i \tau \frac{p-s}{p} \sin \frac{\pi pr's'}{p} \right) \]
\[ \times \left( \frac{r'}{p'} \cos \frac{\pi pr's'}{p' - \tau} - i \tau \frac{p'-s'}{p'} \sin \frac{\pi pr's'}{p'} \right), \quad 1 \leq s \leq p - 1, \]

\[ 1 \leq s' \leq p' - 1, \]

\[ S_{r',s,p'}(\tau) = \frac{\sqrt{2}}{\sqrt{pp'}} \left( -1 \right)^{s+p'+r'p'} \left( \frac{r}{p} \cos \frac{\pi pr's'}{p} - i \tau \frac{p-s}{p} \sin \frac{\pi pr's'}{p} \right), \quad 1 \leq s \leq p - 1 \]

\[ \frac{1}{2} \frac{r'}{pp'}, \quad \text{and the other matrix elements are} \]

\[ \tilde{S}_{r',s,s'}^+(\tau) = (-1)^{rs'+sr'} \frac{\sqrt{2}}{pp'} \left( \frac{pp'rr'}{p^2p'^2} \right) \left( \frac{r}{p} \cos \frac{\pi pr's'}{p} - \frac{\pi pr's'}{p'} \right) \]
\[ + i \tau (p's' - p's) \cos \frac{\pi pr's'}{p} \sin \frac{\pi pr's'}{p'} + i \tau (p's' - p's) \sin \frac{\pi pr's'}{p} \cos \frac{\pi pr's'}{p'} \]
\[ + \left( \frac{p's' - p's}{2} \right)^2 \tau^2 - 2i \pi pp' \tau + \frac{p^2r'^2 + p'^2r^2}{2} \sin \frac{\pi pr's'}{p} \sin \frac{\pi pr's'}{p'}. \]
3.1. The procedure. The steps leading from (2.4) and (2.5) to (1.1), in much the same way as in [9], are as follows.

1. We view the characters in (2.1) as a column vector and write the $S$-transformation formulas as

$$\chi(-\frac{1}{\tau}) = S(\tau)\chi(\tau),$$

with the corresponding $N \times N$ $\tau$-dependent matrix $S(\tau)$, where

$$N = \frac{1}{2}(p-1)(p'-1) + 2pp'.$$

is the total number of characters.

We then take $S(\tau)$ to be the $(2pp') \times (2pp')$ block of $S(\tau)$ corresponding to the $2pp'$ characters $\chi_{r,s}(\tau)$, $1 \leq r \leq p$, $1 \leq s \leq p'$. That is, we deal with only the $S_{r,s}(\tau)$ in (2.6).

In accordance with the block structure of the Jordan form of $S(\tau)$, we fix the block structure of matrices as follows: 2 blocks of size $1 \times 1$, $(p-1)+(p'-1)$ blocks of size $2 \times 2$, and $\frac{1}{2}(p-1)(p'-1)$ blocks of size $4 \times 4$.

2. Totally similarly to [9], there exists a $((2pp') \times (2pp'))$-matric automorphy factor $J(\gamma, \tau)$, with $\gamma \in SL(2, \mathbb{Z})$, satisfying the cocycle condition and a commutativity property formulated in [9], such that

$$S = J(S, \tau)S(\tau)$$

is a numerical ($\tau$-independent) matrix; in fact,

$$S = S(i).$$

(3.1)

It then follows that $S^2 = 1$.

3. From now on, $\chi = (\chi_I)$ denotes the $2pp'$ $W_{p,p'}$-characters ordered as

$$\chi = (\chi_{p,p'}, \chi_{p,p'}, \chi_{p',p}, \chi_{p,p'}, \chi_{p,p'}, \chi_{p',p}, \chi_{p,p'}, \chi_{p',p}, \chi_{p,p'}, \chi_{p',p}, \chi_{p,p'}, \chi_{p',p}, \chi_{p,p'}, \chi_{p',p}, \chi_{p,p'}, \chi_{p',p}).$$

(3.2)

This arrangement of the characters clearly agrees with the $“1+1+(p+p'-2)\cdot 2 \times 2 + \frac{1}{2}(p-1)(p'-1) \cdot 4 \times 4”$ block structure. We also define a special row

$$\widetilde{S}_{r,s}(\tau) = (-1)^{pp'+sp+p}\delta_{r,s}(\tau) - (-1)^{pp'+sp+p}\frac{1}{\sqrt{2pp'}}\sin^{-1}\frac{\pi pp'}{p}\sin^{-1}\frac{\pi pp'}{p'}. $$

In (2.6), remarkably, the dependence on the primed and unprimed indices almost (modulo $(-1)^{pp'+sp'}$) factors, which partly reduces the analysis to that for the $(p,1)$ and $(1,p')$ cases. (Most of the quantum-group objects corresponding to the $(p,p')$ models in [13] also have an “almost factored” form.)

3. "Logarithmic" $(p,p')$-fusion
\[ P_{\Omega} = (1, 1, 1, 0, \ldots, 1, 0, 1, 0, \ldots, 1, 0, 0, 0, \ldots, 1, 0, 0, 0). \]

\[
\begin{align*}
2(p-1) \text{ elements} & \quad 2(p'-1) \text{ elements} & \quad 4{\lfloor \frac{1}{2}(p-1)(p'-1) \rfloor} \text{ elements}
\end{align*}
\]

(At this point, we anticipate that the block structure inherited from \( S(\tau) \) is to become the block structure of the fusion algebra; accordingly, the ones encountered in \( P_{\Omega} \) correspond to the decomposition of the identity into a sum of primitive idempotents, one for each block.)

4. Let \( S_{\Omega} = (S_{\Omega}^J) \) be the row of \( S \) corresponding to the vacuum-representation character \( \chi_{\Omega} = \chi_{1,1}^+ \), i.e.,

\[ \chi_{1,1}^+(-\frac{1}{\tau}) = S_{\Omega}^J \chi_{\Omega}(\tau) \]

(the sum is taken over the \( 2pp' \) values of \( J \) in accordance with (3.2)). With the chosen ordering, \( \chi_{\Omega} \) occupies position \( 2p+2p'-1 \) in (3.2) and, accordingly, \( S_{\Omega} \) is the \( (2p+2p'-1) \)th row. Explicitly (see (2.6)), the segment of \( S_{\Omega} \) corresponding to \( (\chi_{s,s'}, \chi_p, \chi_{s'-s'}, \chi_{p-s', p'-s'}) \) is given by \( \frac{2\sqrt{2}}{pp'p'}(1)^{s'+s} \) times

\[
\begin{align*}
(\cos \frac{\pi ps'}{p} + (p-s) \sin \frac{\pi ps'}{p}) & \left( \cos \frac{\pi ps'}{p'} + (p'-s') \sin \frac{\pi ps'}{p'} \right), \\
(\cos \frac{p s'}{p} - s \sin \frac{p s'}{p}) & \left( \cos \frac{p s'}{p'} + (p'-s') \sin \frac{p s'}{p'} \right), \\
(\cos \frac{p os}{p} + (p-s) \sin \frac{p os}{p}) & \left( \cos \frac{p os}{p'} - s \sin \frac{p os}{p'} \right), \\
(\cos \frac{p os}{p} - s \sin \frac{p os}{p}) & \left( \cos \frac{p os}{p'} - s \sin \frac{p os}{p'} \right).
\end{align*}
\]

5. We next consider the equation (cf. [9])

\[ P_{\Omega} = S_{\Omega}K \]

and solve it for the block-diagonal matrix

\[
K = \begin{pmatrix}
\kappa_1 \\
\kappa_2 \\
K_{2\times 2} \\
K_{2\times 2} \\
K_{4\times 4} \\
\cdots
\end{pmatrix}
\]

The vacuum representation \( \mathcal{K}_{1,1}^+ \) of the \( \mathcal{W}_{p,p'} \) algebra in [12] is in fact an extension of the Virasoro representation \( \chi_{1,1} \) whose character is \( \chi_{1,1}(\tau) \) by the \( \mathcal{W}_{p,p'} \)-representation \( \mathcal{K}_{1,1}^+ \) whose character is \( \chi_{1,1}(\tau) \): 0 \( \rightarrow \) \( \mathcal{K}_{1,1}^+ \rightarrow \mathcal{K}_{1,1}^+ \rightarrow \mathcal{K}_{1,1}^+ \rightarrow 0 \). The difference between \( \mathcal{K}_{1,1}^+ \) and \( \mathcal{K}_{1,1}^+ \) is irrelevant in the present context, where we ignore all the \( \chi_{r,r'} \) characters altogether.
(with zeros outside the blocks), where the $2 \times 2$ blocks are as in [9], i.e., have the structure

$$K_{2\times2} = \begin{pmatrix} a & \lambda \\ -a & b\lambda \end{pmatrix}$$

(it is understood that $K_{2\times2}^{(i)} = \begin{pmatrix} a^{(i)} & \lambda^{(i)} \\ -a^{(i)} & b^{(i)}\lambda^{(i)} \end{pmatrix}$ for each block, but the block dependence is omitted for brevity) and the $4 \times 4$ blocks have the structure

$$K_{4\times4} = \begin{pmatrix} a & \mu & \nu & \frac{1}{a} \mu \nu \\ -a & -\mu & c\nu & \frac{c}{a} \mu \nu \\ -a & b\mu & -\nu & \frac{b}{a} \mu \nu \\ a & -b\mu & -c\nu & \frac{bc}{a} \mu \nu \end{pmatrix}$$

(again, with the block dependence omitted).

The nonzero factors $\lambda$, $\mu$, and $\nu$, rescaling each column except the first in each block, are irrelevant in what follows (because nilpotent elements have no canonical normalization). The unknowns $a$ and $b$ in each $2 \times 2$ block and $a$, $b$, and $c$ in each $4 \times 4$ block are determined from Eq. (3.4). That is, if $(s_1, s_2, s_3, s_4)$ is a segment of $S_{\Omega}$ corresponding to a $4 \times 4$ block, then

$$a = \frac{1}{s_1 - s_2 - s_3 + s_4}, \quad b = \frac{s_2 - s_1}{s_3 - s_4}, \quad c = \frac{s_3 - s_1}{s_2 - s_4}$$

in this block; the equations are compatible because $s_1 s_4 = s_2 s_3$, as is readily seen from (3.3). (By (3.4), the two elements of $K$ that constitute the $1 \times 1$ blocks are the inverse of the corresponding $S$-matrix coefficients, just as the denominators in the semisimple Verlinde formula; with the $2 \times 2$ blocks, the situation repeats that in [9].)

6. We set

$$P = SK.$$

The fusion algebra is now reconstructed from the $P$ matrix in much the same way as in [9], as follows. Evidently, the $(2p + 2p' - 1)$th row of $P$ is just $P_{\Omega}$. Let $P_I$ be the $I$th row of $P$. We define $M_I$, $I = 1, \ldots, 2pp'$, to be block-diagonal matrices that solve the equation

$$P_I = P_{\Omega} M_I$$

and whose $2 \times 2$ blocks are of the form (just as in [9])

$$\begin{pmatrix} \alpha & \beta \\ 0 & \alpha \end{pmatrix}$$
and the $4 \times 4$ blocks are
\[
\begin{pmatrix}
\alpha & \beta & \gamma & \zeta \\
\alpha & 0 & \gamma & \\
\alpha & \beta & & \\
& & & \\
\end{pmatrix}
\]
(with zeros below the diagonal).

The $M_I$ are then determined uniquely; in particular, the $4 \times 4$ blocks are given by
\[
\begin{pmatrix}
p_I & q_I & r_I & s_I \\
p_I & 0 & r_I & \\
p_I & q_I & & \\
p_I & & & \\
\end{pmatrix},
\]
where $(p_I, q_I, r_I, s_I)$ is a segment of $P_I$ corresponding to the chosen block.

The result is then that the $M_I$ satisfy the algebra
\[
M_I M_J = \sum_K n_{IJ}^K M_K,
\]
where the nonnegative integer coefficients $n_{IJ}^K$ turn out to be those read off from (1.1). (Simultaneously, the matrix $N_I = PM_I P^{-1}$ for each $I$ gives the fusion structure constants as $(N_I)_J^K = n_{IJ}^K$.) Evidently, $M_Q$ is the unit in the algebra.

### 3.2. Examples.

The illustrative power of examples is hampered by the rapidly growing matrix size and the general clumsiness of explicit expressions. We consider only the “percolation” and “Lee–Yang” cases, where explicit values of the various matrix entries may be useful for comparison with the studies of these cases by more direct methods (e.g., in [19]).

#### 3.2.1. $(3, 2)$.
For $(p, p') = (3, 2)$, the $12 \times 12$ matrix $S = S(i)$ explicitly evaluates as

|       | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ | $1/2\sqrt{3}$ |
|-------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |
| $1/2\sqrt{3}$ | $1/3$         | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          | $-1/3$         | $1/3$          |

\[S = \begin{pmatrix}
\end{pmatrix} \]
Here, $S_Ω$ is the 9th row of $S$. The matrix $K$ in (3.5) is then given by

\[
K = \begin{pmatrix}
12\sqrt{3} & -12\sqrt{3} & 4 & 1 & -4 & 7-3\sqrt{3} & 4 & 1 & -4 & 7+3\sqrt{3} & -3\sqrt{3} & 1 & 3\sqrt{3} & 1 & -1 & 1 & 1 & 7-3\sqrt{3} & -1 & 1 & -1 & 7-3\sqrt{3} & 3\sqrt{3}-7 & 1 & -1 & 3\sqrt{3}-7 & 3\sqrt{3}-7
\end{pmatrix},
\]

which gives rise to the fusion-algebra eigenmatrix

\[
P = SK = \begin{pmatrix}
6 & -6 & 0 & \frac{3(\sqrt{3}-1)}{2} & 0 & \frac{3(1+2\sqrt{3})}{11} & 0 & \frac{2}{\sqrt{3}} & 0 & 0 & 0 & 6(1-\sqrt{3})
6 & -6 & 0 & \frac{3(\sqrt{3}-1)}{2} & 0 & \frac{3(1+2\sqrt{3})}{11} & 0 & \frac{2}{\sqrt{3}} & 0 & 0 & 0 & 6(\sqrt{3}-1)
2 & -2 & 2 & 0 & -2 & 0 & 0 & \frac{2}{3\sqrt{3}} & 0 & 2 & 0 & 0
4 & -4 & -2 & \frac{3(1-\sqrt{3})}{4} & 2 & -\frac{3(1+2\sqrt{3})}{22} & 0 & \frac{4}{3\sqrt{3}} & 0 & -2 & 0 & 3(\sqrt{3}-1)
4 & -4 & -2 & \frac{3(1-\sqrt{3})}{4} & 2 & -\frac{3(1+2\sqrt{3})}{22} & 0 & \frac{4}{3\sqrt{3}} & 0 & 2 & 0 & 3(1-\sqrt{3})
2 & -2 & 2 & 0 & -2 & 0 & 0 & \frac{2}{3\sqrt{3}} & 0 & -2 & 0 & 0
3 & 3 & 0 & \frac{3(\sqrt{3}-1)}{2} & 0 & -\frac{3(1+2\sqrt{3})}{22} & 3 & 0 & 0 & 0 & 3(1-\sqrt{3}) & 0
3 & 3 & 0 & \frac{3(\sqrt{3}-1)}{2} & 0 & -\frac{3(1+2\sqrt{3})}{22} & -3 & 0 & 0 & 0 & 3(\sqrt{3}-1) & 0
1 & 1 & 1 & 0 & 1 & 0 & 1 & 0 & 1 & 0 & 0 & 0
2 & 2 & -1 & \frac{3(1-\sqrt{3})}{8} & -1 & \frac{3(1+2\sqrt{3})}{44} & 2 & 0 & -1 & 0 & \frac{3(\sqrt{3}-1)}{2} & 0
1 & 1 & 1 & 0 & 1 & 0 & -1 & 0 & -1 & 0 & 0 & 0
2 & 2 & -1 & \frac{3(1-\sqrt{3})}{8} & -1 & \frac{3(1+2\sqrt{3})}{44} & -2 & 0 & 1 & 0 & \frac{3(1-\sqrt{3})}{2} & 0
\end{pmatrix}
\]

The fusion relations that follow in accordance with (3.6)–(3.7) are the $(p = 3, p' = 2)$ specialization of (1.1) (explicitly written in [12]).

3.2.2. $(5, 2)$. For $(p, p') = (5, 2)$, all of the entries of the $20 \times 20$ matrix $S$ can be easily evaluated from the $S_{rr', ss'}(i)$ in (2.6). In particular, the vacuum-representation row is

\[
S_Ω = S_{13} = \begin{pmatrix}
\frac{1}{20\sqrt{3}} & \frac{1}{20\sqrt{3}} & \frac{5-\sqrt{3}+4\sqrt{10(5+\sqrt{5})}}{200} & \frac{5-\sqrt{5}}{200} & 5 \sqrt{3} - 3 \sqrt{10(5+\sqrt{5})} & 5 + \sqrt{5} - 3 \sqrt{10(5-\sqrt{5})}
\frac{5+\sqrt{3}+2\sqrt{10(5-\sqrt{5})}}{200} & 5-\sqrt{5}-2\sqrt{10(5-\sqrt{5})} & -5-\sqrt{3} + 3 \sqrt{10(5-\sqrt{5})} & \sqrt{3} - 5 + \sqrt{10(5+\sqrt{5})} & \sqrt{5} - 5 - 4 \sqrt{10(5+\sqrt{5})} & 5-\sqrt{5} + 4 \sqrt{10(5+\sqrt{5})} & 5-\sqrt{5} - 3 \sqrt{10(5-\sqrt{5})} & \sqrt{5} - 5 - 4 \sqrt{10(5+\sqrt{5})} & \sqrt{5} - 5 + \sqrt{10(5+\sqrt{5})} & 5 + \sqrt{3} - 3 \sqrt{10(5-\sqrt{5})} & \sqrt{5} - 5 + \sqrt{10(5+\sqrt{5})} & 5-\sqrt{5} - 3 \sqrt{10(5-\sqrt{5})} & \sqrt{5} - 5 - 4 \sqrt{10(5+\sqrt{5})}
\end{pmatrix}
\]
The matrix $K$ in (3.5) then consists of the blocks

$$K = \text{diag}\left(20\sqrt{5}, -20\sqrt{5}, \frac{2\sqrt{20(5 - \sqrt{5})}}{20}, \frac{2\sqrt{20(5 + \sqrt{5})}}{20}, \frac{1}{2}, \frac{1}{2}\right),$$

$$\begin{bmatrix}
-2\sqrt{2(5 + \sqrt{5})} & 1 \\
2\sqrt{2(5 + \sqrt{5})} & \frac{109 + 20\sqrt{5} - 5\sqrt{365 + 158\sqrt{5}}}{41}
\end{bmatrix},$$

$$\begin{bmatrix}
\frac{2\sqrt{2(5 - \sqrt{5})}}{20} & 1 \\
-\frac{2\sqrt{2(5 - \sqrt{5})}}{20} & \frac{379 - 40\sqrt{3 - 5\sqrt{6245 - 2558\sqrt{5}}}}{1121}
\end{bmatrix},$$

$$\begin{bmatrix}
\sqrt{\frac{5 + \sqrt{5}}{2}} & 1 \\
-\sqrt{\frac{5 + \sqrt{5}}{2}} & \frac{5\sqrt{5} - 2 + 5\sqrt{5 - 2\sqrt{5}}}{2}
\end{bmatrix}.$$

This gives rise to the fusion-algebra eigenmatrix $P = SK$, shown (at about the limit of reasonable typesetting capabilities) in Fig. 1. The $(p = 5, p' = 2)$-case of algebra (1.1) follows from this $P$ in accordance with (3.6)–(3.7).

### 4. Conclusions

The procedure proposed here is of course not a replacement for the “honest” derivation of fusion (cf. [19]). We also reiterate that the success of this procedure is apparently rooted in the quantum group structure of the corresponding logarithmic conformal field models [12, 13] (and actually amounts to no more than establishing the coincidence with the quantum group Grothendieck ring). For the logarithmic $(p, p')$ models, anyway, the existence of a relation between modular transformations of characters and the fusion additionally supports the “quantum-group candidate” for the fusion of representations of the extended algebra in [12].

But the much more complicated “logarithmic” modular transformations in [21] are not likely to yield a fusion algebra similarly.

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3In fact, Kazhdan–Lusztig-dual quantum groups “know” not only about the numerology and modular group transformations of extended-algebra characters in logarithmic conformal field models but also about the asymptotic form of the characters [20].
The 20 × 20 eigenmatrix \( P = SK \) for \((p, p') = (5, 2)\). The vacuum-representation row is the 13th.
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