SKEW-UNFOLDING THE SKOROKHOD REFLECTION OF A CONTINUOUS SEMIMARTINGALE

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Dedicated to Terry Lyons on the occasion of his 60th birthday

Abstract

The Skorokhod reflection of a continuous semimartingale is unfolded, in a possibly skewed manner, into another continuous semimartingale on an enlarged probability space according to the excursion-theoretic methodology of Prokaj (2009). This is done in terms of a skew version of the Tanaka equation, whose properties are studied in some detail. The result is used to construct a system of two diffusive particles with rank-based characteristics and skew-elastic collisions. Unfoldings of conventional reflections are also discussed, as are examples involving skew Brownian Motions and skew Bessel processes.

Keywords and Phrases: Skorokhod and conventional reflections; skew and perturbed Tanaka equations; skew Brownian and Bessel processes; pure and Ocone martingales; local time; competing particle systems; asymmetric collisions.

AMS 2000 Subject Classifications: Primary, 60G42; secondary, 60H10.

1 The Result

On a filtered probability space \((\Omega, \mathcal{F}, \mathbb{P})\), \(\mathcal{F} = \{\mathcal{F}(t)\}_{0 \leq t < \infty}\) satisfying the so-called “usual conditions” of right continuity and augmentation by null sets, we consider a real-valued continuous semimartingale \(U(\cdot)\) of the form

\[
U(t) = M(t) + A(t), \quad 0 \leq t < \infty
\]

(1.1)

with \(M(\cdot)\) a continuous local martingale and \(A(\cdot)\) a process of finite first variation on compact intervals. We assume \(M(0) = A(0) = 0\) for concreteness.

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There are two ways to “fold”, or reflect, this semimartingale about the origin. One is the \textit{conventional reflection}
\begin{equation}
R(t) := |U(t)|, \quad 0 \leq t < \infty; \tag{1.2}
\end{equation}
the other is the \textit{SKOROKHOD reflection}
\begin{equation}
S(t) := U(t) + \max_{0 \leq s \leq t} (-U(s)), \quad 0 \leq t < \infty. \tag{1.3}
\end{equation}

The following result, inspired by PROKAJ (2009), shows how the first can be obtained from the second, by suitably unfolding the \textit{SKOROKHOD reflection} in a possibly “skewed” manner.

**Theorem 1.1.** Fix a constant \(\alpha \in (0, 1)\). There exists an enlargement \((\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})\), \(\tilde{\mathcal{F}} = \{\tilde{\mathcal{F}}(t)\}_{0 \leq t < \infty}\) of the filtered probability space \((\Omega, \mathcal{F}, \mathbb{P})\), \(\mathcal{F} = \{\mathcal{F}(t)\}_{0 \leq t < \infty}\) with a measure-preserving map \(\pi : \Omega \to \tilde{\Omega}\), and on this enlarged space a continuous semimartingale \(X(\cdot)\) that satisfies
\begin{equation}
|X(\cdot)| = S(\cdot), \quad L^X(\cdot) = \alpha L^S(\cdot), \quad X(\cdot) = \int_0^\cdot \text{sgn}(X(t)) \, dU(t) + \frac{2 \alpha - 1}{\alpha} L^X(\cdot). \tag{1.4}
\end{equation}

Here and throughout this note, we use the notation
\begin{equation}
L^U(\cdot) := \lim_{\varepsilon \downarrow 0} \frac{1}{2 \varepsilon} \int_0^\cdot 1_{\{0 \leq U(t) < \varepsilon\}} \, d(U(t)), \quad \hat{L}^U(\cdot) := \frac{1}{2} \left( L^U(\cdot) + L^{-U}(\cdot) \right) \tag{1.5}
\end{equation}
respectively for the \textit{right} and the \textit{symmetric} local time at the origin of a continuous semimartingale as in \((1.1)\), and the conventions
\[\text{sgn}(x) := 1_{(0, \infty)}(x) - 1_{(-\infty, 0)}(x), \quad \text{sgn}(x) := 1_{(0, \infty)}(x) - 1_{(-\infty, 0)}(x), \quad x \in \mathbb{R}\]
for the symmetric and the left-continuous versions, respectively, of the signum function. We also denote by \(\mathbb{F}^U = \{\mathcal{F}^U(t)\}_{0 \leq t < \infty}\) the “natural filtration” of \(U(\cdot)\), that is, the smallest filtration that satisfies the usual conditions and with respect to which \(U(\cdot)\) is adapted; we set \(\mathcal{F}^U(\infty) := \sigma\left(\bigcup_{0 \leq t < \infty} \mathcal{F}^U(t)\right)\). Equalities between stochastic processes, such as in \((1.4)\), are to be understood throughout in the almost sure sense.

Theorem\(\ref{thm:1.1}\) constructs a continuous semimartingale \(X(\cdot)\) whose conventional reflection coincides with the SKOROKHOD reflection of the given semimartingale \(U(\cdot)\), and which satisfies the stochastic integral equation in \((1.4)\). We think of this equation as a \textit{skew version} of the celebrated TANAKA equation driven by the continuous semimartingale \(U(\cdot)\), whose “skew-unfolding” it produces via the parameter \(\alpha\). When there is no skewness, i.e., with \(\alpha = 1/2\), the integral equation of \((1.4)\) reduces to the classical TANAKA equation; in this case Theorem\(\ref{thm:1.1}\) is just the main result in the paper by PROKAJ (2009), which inspired our work.

We shall prove Theorem\(\ref{thm:1.1}\) in Section\(\ref{sec:4}\) then use it in Section\(\ref{sec:5}\) to construct a system of two diffusive particles with rank-based characteristics and skew-elastic collisions. Section\(\ref{sec:4}\) discusses a similar skew-unfolding of the conventional reflection \(R(\cdot) = |U(\cdot)|\) of \(U(\cdot)\). In the section that follows we discuss briefly some properties of the \textit{skew TANAKA equation} in \((1.4)\).

## 2 The Skew TANAKA Equation

A first question that arises regarding the stochastic integral equation in \((1.4)\), is whether it can be written in the more conventional form
\begin{equation}
X(\cdot) = \int_0^\cdot \text{sgn}(X(t)) \, dU(t) + \frac{2 \alpha - 1}{\alpha} L^X(\cdot), \tag{2.1}
\end{equation}
in terms of the asymmetric (left-continuous) version of the signum function.
For this, it is necessary and sufficient to have
\[
\int_0^\infty 1_{\{X(t) = 0\}} \, dU(t) \equiv 0, \quad \text{or equivalently} \quad \int_0^\infty 1_{\{S(t) = 0\}} \, dU(t) \equiv 0 \tag{2.2}
\]
in the context of Theorem 1.1. Now from (1.1), (1.3) it is clear that \( M(\cdot) \) is the local martingale part of the continuous semimartingale \( S(\cdot) \), so we have \( \langle S(\cdot) \rangle = \langle U(\cdot) \rangle = \langle M(\cdot) \rangle \) and
\[
\int_0^\infty 1_{\{S(t) = 0\}} \, d\langle M(\cdot) \rangle(t) = 0 \tag{2.3}
\]
(e.g., Karatzas & Shreve, Exercise 3.7.10). This gives \( \int_0^\infty 1_{\{S(t) = 0\}} \, dM(t) \equiv 0 \), so (2.2) will follow if and only if
\[
\int_0^\infty 1_{\{S(t) = 0\}} \, dA(t) \equiv 0 \tag{2.4}
\]
holds; and on the strength of (2.3), a sufficient condition for (2.4) is that there exist an \( \langle u \rangle \)-progressively measurable process \( p(\cdot) \), locally integrable with respect to \( \langle M(\cdot) \rangle \) and such that
\[
A(\cdot) = \int_0^\cdot p(t) \, d\langle M(\cdot) \rangle(t) \tag{2.6}
\]
Proof: The first and third claims have already been argued. As for the second, we observe that the Itô-Tanaka formula applied to (2.1) gives
\[
S(\cdot) = |X(\cdot)| = \int_0^\cdot \text{sgn}(X(t)) \, dX(t) + 2L^X(\cdot) = U(\cdot) - \frac{2\alpha - 1}{\alpha} L^X(\cdot) + 2L^X(\cdot) = U(\cdot) + L^S(\cdot)
\]
on the strength of the second equality in (1.4). It is clear from this expression that the filtration comparison \( \mathcal{F}^U(t) \subseteq \mathcal{F}^S(t) \) holds for all \( 0 \leq t < \infty \); whereas the reverse inclusion and the claimed identification are direct consequences of (1.3). \( \square \)
Remark: More generally (that is, in the absence of condition (2.4)), the local time at the origin of the Skorokhod reflection \( S(\cdot) \) is \( L^S(t) = \max_{0 \leq s \leq t} (-U(s)) + \int_0^t 1_{\{S(u) = 0\}} \, dA(u), \quad 0 \leq t < \infty. \)

2.1 Uniqueness in Distribution for the Skew Tanaka Equation

A second question that arises regarding the skew-Tanaka equation of (1.4), is whether it can be solved uniquely. It is well-known that we cannot expect pathwise uniqueness or strength to hold for this equation. Such strong existence and uniqueness fail already with \( \alpha = 1/2 \) and \( U(\cdot) \) a standard Brownian motion, in which case we have in (2.5) also the strict inclusion \( \mathcal{F}^U(t) \subsetneq \mathcal{F}^S(t) \) for all \( t \in (0, \infty) \) (e.g., Karatzas & Shreve (1991), Example 5.3.5). The Skorokhod reflection of \( U(\cdot) \) can then be “unfolded” into a Brownian motion \( X(\cdot) \), whose filtration is strictly finer than that of the original Brownian motion \( U(\cdot) \); the unfolding cannot be accomplished without the help of some additional randomness.
The issue, therefore, is whether uniqueness in distribution holds for the skew-TANAKA equation of (1.4), under appropriate conditions. We shall address this question in the case of a continuous local martingale \( U(\cdot) \) with \( U(0) = 0 \) and \( \langle U \rangle(\infty) = \infty \). Let us recall a few notions and facts about such a process, starting with its DAMBIS-DUBINS-SCHWARZ representation
\[
U(t) = B(\langle U \rangle(t)), \quad 0 \leq t < \infty
\]
(cf. KARATZAS & SHREVE (1991), Theorem 3.4.6); here \( B(\theta) = U(Q(\theta)), \, 0 \leq \theta < \infty \) is standard Brownian motion, and \( Q(\cdot) \) the right-continuous inverse of the continuous, increasing process \( \langle U \rangle(\cdot) \).

We say that this \( U(\cdot) \) is pure, if each \( \langle U \rangle(t) \) is \( \mathcal{F}_t^\omega \)-measurable; we say that it is an Ocone martingale, if the processes \( B(\cdot) \) and \( \langle U \rangle(\cdot) \) are independent (cf. Ocone (1993) and Dubins et al. (1993), Appendix). As discussed in Vostrikova & Yor (2000), a pure Ocone martingale is a Gaussian process.

**Proposition 2.2.** Suppose that \( U(\cdot) \) is a continuous local martingale with \( U(0) = 0 \) and \( \langle U \rangle(\infty) = \infty \). Then uniqueness in distribution holds for the skew-TANAKA equation of (1.4), or equivalently of (2.7), provided that either

(i) \( U(\cdot) \) is pure; or that

(ii) the quadratic variation process \( \langle U \rangle(\cdot) \) is adapted to a Brownian motion \( \Gamma(\cdot) := (\Gamma_1(\cdot), \ldots, \Gamma_n(\cdot)), \) with values in some Euclidean space \( \mathbb{R}^n \) and independent of the real-valued Brownian motion \( B(\cdot) \) in the representation (2.7).

**Proof:** Let us consider a continuous local martingale \( U(\cdot) \) with \( U(0) = 0 \), and any continuous semimartingale \( X(\cdot) \) that satisfies the stochastic integral equation in (1.4). Then \( X(\cdot) \) also satisfies the equation of (2.1), as the condition (2.6) holds in this case trivially with \( p(\cdot) \equiv 0 \). In fact, the equation (2.1) can be written then in the form
\[
X(Q(s)) = \int_0^s \text{sgn}(X(Q(\theta))) \, dB(\theta) + \frac{2\alpha - 1}{\alpha} L^X(Q(s)), \quad 0 \leq s < \infty,
\]
with \( Q(\cdot) \) the right-continuous inverse of the continuous, increasing process \( \langle U \rangle(\cdot) \); cf. Proposition 3.4.8 in Karatzas & Shreve (1991). Setting
\[
\tilde{X}(s) := X(Q(s)), \quad \text{it is straightforward to check} \quad L^\tilde{X}(s) = L^X(Q(s)), \quad 0 \leq s < \infty;
\]
for this, one uses the representation (1.5) for the local time at the origin, along with the fact that the local martingale part of the continuous semimartingale \( X(\cdot) \) in (2.1) has quadratic variation process \( \langle U \rangle(\cdot) \). Thus, the time-changed process \( \tilde{X}(\cdot) \) satisfies the stochastic integral equation
\[
\tilde{X}(s) = \int_0^s \text{sgn}(\tilde{X}(\theta)) \, dB(\theta) + \frac{2\alpha - 1}{\alpha} L^\tilde{X}(s), \quad 0 \leq s < \infty.
\]

This can be cast as the Harrison-Shepp (1981) equation
\[
\tilde{X}(\cdot) = \tilde{W}(\cdot) + \frac{2\alpha - 1}{\alpha} L^\tilde{X}(\cdot)
\]
for the skew Brownian motion, driven by the standard Brownian motion
\[
\tilde{W}(\cdot) := \int_0^\cdot \text{sgn}(\tilde{X}(\theta)) \, dB(\theta).
\]
It is well-known from the theory of HARRISON & SHEPP (1981) that the equation (2.9) has a pathwise unique, strong solution; in fact, the skew Brownian motion $\tilde{X}(\cdot)$ and the Brownian motion $\tilde{W}(\cdot)$ generate the same filtration. Since
\[
X(t) = \tilde{X}(\langle U \rangle(t)), \quad 0 \leq t < \infty
\]
holds with $\tilde{X}(\cdot)$ adapted to $\mathbb{P}^{\tilde{W}}$, the distribution of $X(\cdot)$ is uniquely determined whenever
\[
\text{the Brownian motion } \tilde{W}(\cdot) \text{ of (2.10) is independent of the process } \langle U \rangle(\cdot),
\]
or whenever
\[
\langle U \rangle(t) \text{ is } \mathcal{F}^{\tilde{W}}(\infty) - \text{measurable, for every } t \in [0, \infty).
\]

But (2.13) holds when $U(\cdot)$ is pure (case (i) of the Proposition); this is because from (2.10) we have $B(\cdot) = \int_0^\cdot \text{sgn}(\tilde{X}(\theta)) \, d\tilde{W}(\theta)$, therefore $\mathcal{F}^B(t) \subseteq \mathcal{F}^{\tilde{W}}(t)$ for all $t \in [0, \infty)$ and thus $\mathcal{F}^B(\infty) \subseteq \mathcal{F}^{\tilde{W}}(\infty)$.

On the other hand, (2.12) holds under the condition of case (ii) in the Proposition, as $\langle U \rangle(\cdot)$ is then adapted to the filtration generated by the $n$-dimensional Brownian motion $\Gamma(\cdot)$; this, in turn, is independent of $\tilde{W}(\cdot)$ on the strength of the P. Lévy Theorem (e.g., KARATZAS & SHREVE, Theorem 3.3.16), since
\[
\langle \tilde{W}, \Gamma_j \rangle(\cdot) = \int_0^\cdot \text{sgn}(\tilde{X}(\theta)) \, dB, \Gamma_j(\theta) \equiv 0, \quad \forall \ j = 1, \cdots, n.
\]
The proof of the proposition is complete.

**Remark 2.1.** It would be interesting to obtain sufficient conditions for either (2.12) or (2.13) to hold, which are weaker than those of Proposition 2.2. As Example 2.2 shows, however – and contrary to our own initial guess – we cannot expect the conclusions of Proposition 2.2 to remain true for general Ocone martingales.

**Example 2.1. From Brownian Motion to Skew Brownian Motion:** Suppose that $U(\cdot)$ is standard, real valued Brownian motion. Then the conditions of Propositions 2.1 and 2.2 are satisfied rather trivially; uniqueness in distribution holds for the skew-TANAKA equation of (2.1) (equivalently, of (1.4)); and every continuous semimartingale $X(\cdot)$ that satisfies (2.1) is of the form
\[
X(\cdot) = W(\cdot) + \frac{2\alpha - 1}{\alpha} \overset{\text{symmetric}}{L} X(\cdot) \quad \text{with} \quad W(\cdot) := \int_0^\cdot \text{sgn}(X(t)) \, dU(t),
\]
or equivalently
\[
X(\cdot) = W(\cdot) + 2(2\alpha - 1) \overset{\text{symmetric}}{L} X(\cdot)
\]
in terms of the symmetric local time as in (1.5). Of course $W(\cdot)$ is standard Brownian motion by the P. Lévy theorem, and the HARRISON-SHEPP (1981) theory once again characterizes $X(\cdot)$ as skew Brownian motion with parameter $\alpha$. The processes $W(\cdot)$ and $X(\cdot)$ generate the same filtration, which is strictly finer than the filtration generated by the original Brownian motion $U(\cdot) = \int_0^\cdot \text{sgn}(X(t)) \, dW(t)$.

**Example 2.2. Failure of Uniqueness in Distribution for General Ocone Martingales:** We adapt to our setting a construction from page 131 of DUBINS ET AL. (1993). We start with a filtered probability space $(\Omega, \mathcal{F}, \mathbb{P})$, $\mathbb{P}^B = \{\mathcal{F}^B(t)\}_{0 \leq t < \infty}$ where $B(\cdot)$ is standard Brownian motion with $B(0) = 0$, and define the adapted, continuous and strictly increasing process
\[
A(t) := t \cdot 1_{\{t \leq 1\}} + \left\{1 + (u \cdot 1_{\{B(0) > 0\}} + v \cdot 1_{\{B(0) \leq 0\}})(t - 1)\right\} \cdot 1_{\{t > 1\}}, \quad 0 \leq t < \infty
\]
where $u > 0$ and $v > 0$ are given real numbers with $u \neq v$, as well as the processes
\[
X(\cdot) := B(A(\cdot)), \quad \Xi(\cdot) := -X(\cdot).
\]
The Lévy transform

\[ \beta(\cdot) := \int_0^\cdot \text{sgn}(B(t)) \, dB(t) \]

of \( B(\cdot) \) is a standard Brownian motion adapted to the filtration \( \mathbb{F}^B = \{ \mathcal{F}^B_t \}_{0 \leq t < \infty} \), which is strictly coarser than \( \mathbb{F}^B \); in particular, it can be seen that \( \beta(\cdot) \) is independent of \( \text{sgn}(B(1)) = 2 \mathbf{1}_{\{B(1) > 0\}} - 1 \), and thus of the process \( A(\cdot) \) as well.

On the other hand, the process \( X(\cdot) \) is a martingale of its natural filtration \( \mathbb{F}^X = \{ \mathcal{F}^X_t \}_{0 \leq t < \infty} \); therefore, so is its “mirror image” \( \Xi(\cdot) \), and more importantly its Lévy transform

\[ U(\cdot) := \int_0^\cdot \text{sgn}(X(t)) \, dX(t) = \beta(A(\cdot)) \quad \text{with} \quad \langle U \rangle(\cdot) = A(\cdot), \]

which is thus seen to be an Ocone martingale. Now clearly, both \( X(\cdot) \) and \( \Xi(\cdot) \) satisfy the equation (2.11) with \( \alpha = 1/2 \) driven by \( U(\cdot) \), so pathwise uniqueness fails for this equation. We also note that the conditions of Proposition 2.2 fail too in this case.

We claim that uniqueness in distribution fails as well. In a manner similar to the treatment in Dubins et al. (1993), we shall argue that the distributions of \( X(\cdot) \) and \( \Xi(\cdot) \) at time \( t = 2 \) are different. Now if the random variables

\[ X(2) = B(1 + u) \cdot 1_{\{B(1) > 0\}} + B(1 + v) \cdot 1_{\{B(1) \leq 0\}} \quad \text{and} \quad \Xi(2) = -X(2) \]

had the same probability distributions, that is, if the distribution of the random variable \( X(2) \) were symmetric about the origin, we would have \( \mathbb{E}[\langle X(2) \rangle^3] = 0 \). However, let us note the decomposition

\[ X(2) = B(1) + (B(1 + u) - B(1)) \cdot 1_{\{B(1) > 0\}} + (B(1 + v) - B(1)) \cdot 1_{\{B(1) \leq 0\}}, \]

which gives

\[ \mathbb{E}[\langle X(2) \rangle^3] = 3 \mathbb{E} \left[ B(1) \left( B(1 + u) - B(1) \right)^2 1_{\{B(1) > 0\}} \right] + 3 \mathbb{E} \left[ B(1) \left( B(1 + v) - B(1) \right)^2 1_{\{B(1) \leq 0\}} \right] \]

\[ = 3 \mathbb{E} \left[ (B(1))^+ \right] (u - v) \neq 0. \]

This contradiction establishes the claim.

2.2 The Perturbed Skew-Tanaka Equation is Strongly Solvable

The addition of some independent noise can restore pathwise uniqueness, thus also strength, to weak solutions of the stochastic equation in (1.4) or (2.1). In the spirit of Prokaj (2013) or Fernholz, Ichiba, Karatzas & Prokaj (2013), hereafter referred to as [FIKP], we have the following result.

Proposition 2.3. Suppose that the continuous semimartingale \( U(\cdot) \) as in (1.1) satisfies the conditions of Proposition 2.1 where now the \( \mathbb{F} \)-progressively measurable process \( p(\cdot) \) of (2.6) is locally square-integrable with respect to \( \langle M \rangle(\cdot) \); and that

\[ V(\cdot) = N(\cdot) + \Delta(\cdot) \]

is another continuous semimartingale, with continuous local martingale part \( N(\cdot) \) and finite variation part \( \Delta(\cdot) \) which satisfy \( N(0) = \Delta(0) = 0 \) and

\[ \langle M, N \rangle(\cdot) \equiv 0, \quad \langle M \rangle(\cdot) = \int_0^\cdot q(t) \, d\langle N \rangle(t) \]

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for some $\mathbb{F}$–progressively measurable process $q(\cdot)$ with values in a compact interval $[0, b]$.

Then pathwise uniqueness holds for the perturbed skew-TANAKA equation

$$
X(\cdot) = \int_0^\cdot \text{sgn}(X(t)) \, dU(t) + V(\cdot) + \frac{2\alpha - 1}{\alpha} L X(\cdot),
$$

provided that either

(i) $\alpha = 1/2$, or that

(ii) $U(\cdot)$ and $V(\cdot)$ are independent, standard Brownian motions. In this case a weak solution to (2.16) exists, and is thus strong.

The claim of case (i) is proved in Theorem 8.1 of [FIKP], and the claim of case (ii) in an Appendix, section 6. In case (ii) of Proposition 2.3 the equation (2.16) can be written equivalently as

$$
X(\cdot) = \int_0^\cdot 1_{\{X(t)>0\}} \, dW_+(t) + \int_0^\cdot 1_{\{X(t)<0\}} \, dW_-(t) + \frac{2\alpha - 1}{\alpha} L X(\cdot).
$$

Here $W_\pm(\cdot) := V(\cdot) \pm U(\cdot)$ are independent Brownian motions with local variance 2; one of them governs the motion of $X(\cdot)$ during its positive excursions, the other during the negative ones, whereas these excursions get skewed when $\alpha \neq 1/2$.

3 Proof of Theorem 1.1

We shall follow very closely the methodology of PROKAI (2009), with some necessary modifications related to the skewness. The enlargement of the filtered probability space $(\Omega, \mathcal{F}, \mathbb{P})$, $\mathbb{P} = \{\mathcal{F}(t)\}_{0 \leq t < \infty}$ is done in terms of a sequence $\{\xi_k\}_{k \in \mathbb{N}}$ of independent random variables with common BERNOULLI distribution

$$
\mathbb{P}(\xi_1 = +1) = \alpha, \quad \mathbb{P}(\xi_1 = -1) = 1 - \alpha
$$

(3.1)

(thus with expectation $\mathbb{E}(\xi_1) = 2\alpha - 1$), which is independent of $\mathcal{F}(\infty) = \sigma(\bigcup_{0 \leq t < \infty} \mathcal{F}(t))$. On the enlarged probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$ we have all the objects of the original space, so we keep the same notation for them. We denote by

$$
\mathcal{Z} := \{t \geq 0 : S(t) = 0\}
$$

(3.2)

the zero set of the SKOROKHOD reflection $S(\cdot)$ in (1.3), and enumerate as $\{C_k\}_{k \in \mathbb{N}}$ the disjoint components of $[0, \infty) \setminus \mathcal{Z}$, that is, the countably-many excursion intervals of the process $S(\cdot)$ away from the origin. This we do in a measurable manner, so that

$$
\{t \in C_k\} \in \mathcal{F}(\infty), \quad \forall \ t \geq 0, \ k \in \mathbb{N}.
$$

In order to simplify notation, we set

$$
C_0 := \mathcal{Z}, \quad \xi_0 := 0.
$$

(3.3)

We define now

$$
Z(t) := \sum_{k \in \mathbb{N}_0} \xi_k \mathbf{1}_{C_k}(t), \quad \tilde{\mathcal{F}}(t) := \mathcal{F}(t) \lor \mathcal{F}^Z(t)
$$

(3.4)

for all $t \in [0, \infty)$; this gives the enlarged filtration $\tilde{\mathbb{F}} = \{\tilde{\mathcal{F}}(t)\}_{0 \leq t < \infty}$. We posit the following two claims.

**Proposition 3.1.** The process $M(\cdot)$ of (1.7) is a continuous local martingale of the enlarged filtration $\tilde{\mathbb{F}}$. Consequently, both $U(\cdot)$ and $S(\cdot)$ are continuous $\tilde{\mathbb{F}}$–semimartingales.
Proposition 3.2. In the notation of (1.3) and (3.4), we have
\[ Z(\cdot) S(\cdot) = \int_0^\cdot Z(t) \, dS(t) + (2\alpha - 1) \, L^S(\cdot). \]  

Taking the claims of these two propositions at face-value for a moment, we can proceed with the proof of Theorem 1.1 as follows. We define the process
\[ X(\cdot) := Z(\cdot) S(\cdot) \]  
and note
\[ Z(\cdot) = \text{sgn}(X(\cdot)) \quad |X(\cdot)| = S(\cdot) \]  
thanks to (3.3) and (3.4), as well as
\[ X(\cdot) - \int_0^\cdot \text{sgn}(X(t)) \, dS(t) = Z(\cdot) S(\cdot) - \int_0^\cdot Z(t) \, dS(t) = (2\alpha - 1) \, L^S(\cdot) \]  
thanks to (3.6) and (3.5). In particular, \( X(\cdot) \) is an \( \wt{F} \)-semimartingale, and we note the property
\[ 2 \, L^X(\cdot) - L^S(\cdot) = 2 \, L^X(\cdot) - L^{|X|}(\cdot) = \int_0^\cdot \mathbb{1}_{\{X(t) = 0\}} \, dX(t) \]  
of its local time at the origin (cf. section 2.1 in Ichiba et al. (2013)). In conjunction with (3.8) and the fact that \( X(\cdot), S(\cdot), \) and \( Z(\cdot) \) all have the same zero set \( \mathcal{Z} \) as in (3.2), (3.3), we get from this last equation
\[ 2 \, L^X(\cdot) - L^S(\cdot) = \int_0^\cdot \mathbb{1}_{\{X(t) = 0\}} \left[ \text{sgn}(X(t)) \, dS(t) + (2\alpha - 1) \, L^S(t) \right] = (2\alpha - 1) \, L^S(\cdot), \]  
thus
\[ L^X(\cdot) = \alpha \, L^S(\cdot), \]  
establishing the second equality in (1.4). Back in (3.8), this leads to
\[ X(\cdot) = \int_0^\cdot \text{sgn}(X(t)) \left[ dU(t) + dC(t) \right] + (2\alpha - 1) \, L^S(\cdot), \]  
where \( C(\cdot) \) is the continuous, adapted and increasing process
\[ C(t) := S(t) - U(t) = \max_{0 \leq s \leq t} \left( -U(s) \right), \quad 0 \leq t < \infty. \]  

From the theory of the Skorokhod reflection problem we know that this process \( C(\cdot) \) is flat off the set \( \{ t \geq 0 : S(t) = 0 \} = \mathcal{Z} \), so the skew-Tanaka equation of (1.4) follows now from (3.11) and (3.10).

The proof of Theorem 1.1 is complete.

Proof of Proposition 3.1. By localization of necessary, it suffices to show that if \( M(\cdot) \) is an \( \mathbb{F} \)-martingale, then it is also an \( \wt{F} \)-martingale; that is, for any given \( 0 < \theta < t < \infty \) and \( A \in \mathcal{F}(\theta) \) we have
\[ \mathbb{E} \left[ \left( M(t) - M(\theta) \right) \mathbb{1}_A \right] = 0. \]  

It is clear from (3.4) that we need to consider only sets of the form \( A = B \cap D \), where \( B \in \mathcal{F}(\theta) \) and
\[ D = \bigcap_{j=1}^n \{ Z(t_j) = \varepsilon_j \} = \bigcap_{j=1}^n \{ \xi_k(t_j) = \varepsilon_j \} \]  

(3.13)
for \( n \in \mathbb{N} \), \( 0 < t_1 < t_2 < \cdots < t_n < \theta < t \) and \( \varepsilon \in \{-1, 0, 1\} \). Here we have denoted by \( \kappa(u) \) the (random) index of the excursion interval \( C_k \) to which a given \( u \in [0, \infty) \) belongs.

For such choices, and because

\[
\mathbb{E} \left[ (M(t) - M(\theta)) \mathbf{1}_A \right] = \mathbb{E} \left[ (M(t) - M(\theta)) \mathbf{1}_B \cdot \mathbb{E} \left( \mathbf{1}_D \mid \mathcal{F}(\infty) \right) \right],
\]

we see that, in order to prove (3.12), it is enough to argue that

\[
\mathbb{E} \left( \mathbf{1}_D \mid \mathcal{F}(\infty) \right) \text{ is } \mathcal{F}(\theta) \text{– measurable.} \tag{3.14}
\]

But the random variables \( \kappa(t_j) \) in (3.13) are measurable with respect to \( \mathcal{F}(\infty) \), whereas the random variables \( \xi_1, \xi_2, \cdots \) are independent of this \( \sigma \)–algebra. Therefore, we have

\[
\mathbb{E} \left( \mathbf{1}_D \mid \mathcal{F}(\infty) \right) = \mathbb{P} \left[ \bigcap_{j=1}^n \{ \xi_{\kappa(t_j)} = \varepsilon_j \} \mid \mathcal{F}(\infty) \right] = \mathbb{P} \left( \xi_{k_1} = \varepsilon_1, \cdots, \xi_{k_n} = \varepsilon_n \right) \bigg|_{k_1 = \kappa(t_1), \cdots, k_n = \kappa(t_n)}.
\]

For given indices \( (k_1, \cdots, k_n) \) and \( (\varepsilon_1, \cdots, \varepsilon_n) \), let us denote by \( m \) the number of distinct non-zero indices in \( (k_1, \cdots, k_n) \), by \( \lambda \) the number from among those distinct indices of the corresponding \( \varepsilon_j \)’s that are equal to 1, and observe

\[
\mathbb{P} \left( \xi_{k_1} = \varepsilon_1, \cdots, \xi_{k_n} = \varepsilon_n \right) = 0, \text{ if } (\varepsilon_1, \cdots, \varepsilon_n) \text{ contradicts } (k_1, \cdots, k_n); \]

\[
= \alpha^\lambda \left( 1 - \alpha \right)^{m-\lambda}, \text{ otherwise.} \tag{3.16}
\]

Here “\( (\varepsilon_1, \cdots, \varepsilon_n) \) contradicts \( (k_1, \cdots, k_n) \)” means that we have either

. \( k_i = k_j \) but \( \varepsilon_i \neq \varepsilon_j \) for some \( i \neq j \); or
. \( k_i = 0 \) but \( \varepsilon_i \neq 0 \), for some \( i \); or
. \( k_i \neq 0 \) but \( \varepsilon_i = 0 \), for some \( i \).

We note now that when \( k_1 = \kappa(t_1), \cdots, k_n = \kappa(t_n) \), the value of \( m \) (that is, the number of excursion intervals in \([0, s]\) that contain some \( t_i \)), the value of \( \lambda \) (i.e., the number of such excursion intervals that are positive) and the statement “\( (\varepsilon_1, \cdots, \varepsilon_n) \) contradicts \( (k_1, \cdots, k_n) \)”, can all be determined on the basis of the trajectory \( S(u), 0 \leq u \leq \theta \). That is, the quantity on the right-hand side of (3.15) is \( \mathcal{F}^S(\theta) \) – measurable. As a consequence, the property (3.14) holds.

**Proof of Proposition 3.2** For any \( \varepsilon \in (0, 1) \) we define recursively, starting with \( \tau^\varepsilon_0 := 0 \), a sequence of stopping times

\[
\tau^\varepsilon_{2\ell+1} := \inf \left\{ t > \tau^\varepsilon_{2\ell} : S(t) > \varepsilon \right\}, \quad \tau^\varepsilon_{2\ell+2} := \inf \left\{ t > \tau^\varepsilon_{2\ell+1} : S(t) = 0 \right\}
\]

for \( \ell \in \mathbb{N}_0 \). We use this sequence to approximate the process \( Z(\cdot) \) of (3.4) by

\[
Z^\varepsilon(t) := \sum_{\ell \in \mathbb{N}_0} Z(t) \mathbf{1}_{(\tau^\varepsilon_{2\ell+1}, \tau^\varepsilon_{2\ell+2})}(t), \quad 0 \leq t < \infty.
\]

Let us note that the resulting process \( Z^\varepsilon(\cdot) \) is constant on each of the indicated intervals; that the sequence of stopping times just defined does not accumulate on any bounded time-interval, on account of the fact that \( S(\cdot) \) has continuous paths; and that the process \( Z^\varepsilon(\cdot) \) is of finite first variation over compact intervals. We deduce

\[
Z^\varepsilon(T) S(T) = \int_0^T Z^\varepsilon(t) dS(t) + \int_0^T S(t) dZ^\varepsilon(t), \quad 0 \leq T < \infty. \tag{3.17}
\]
The piecewise-constant process \( Z^\varepsilon(\cdot) \) tends to \( Z(\cdot) \) pointwise as \( \varepsilon \downarrow 0 \), and we have

\[
\lim_{\varepsilon \downarrow 0} \int_0^T Z^\varepsilon(t) \, dS(t) = \int_0^T Z(t) \, dS(t), \quad \text{in probability}
\]

(3.18)

for any given \( T \in [0, \infty) \); all the while, \( |Z^\varepsilon(\cdot)| \leq 1 \). On the other hand, the second integral in (3.17) can be written as

\[
\int_0^T S(t) \, dZ^\varepsilon(t) = \sum_{\{\ell : \tau_{2\ell+1}^\varepsilon < T\}} S(\tau_{2\ell+1}^\varepsilon) Z(\tau_{2\ell+1}^\varepsilon) = \varepsilon \sum_{\{\ell : \tau_{2\ell+1}^\varepsilon < T\}} Z(\tau_{2\ell+1}^\varepsilon)
\]

\[
= \varepsilon \sum_{j=1}^{N(T, \varepsilon)} \xi_{\ell_j} = \varepsilon N(T, \varepsilon) \cdot \frac{1}{N(T, \varepsilon)} \sum_{j=1}^{N(T, \varepsilon)} \xi_{\ell_j},
\]

where \( \{\xi_{\ell_j}\}_{j=1}^{N(T, \varepsilon)} \) is an enumeration of the values \( Z(\tau_{2\ell+1}^\varepsilon) \) and

\[
N(T, \varepsilon) := \# \{\ell : \tau_{2\ell+1}^\varepsilon < T\}
\]

is the number of upcrossings of the interval \((0, \varepsilon)\) that the process \( S(\cdot) \) has completed by time \( T \). From Theorem VI.1.10 in Revuz & Yor (1999), we have the representation of local time \( \lim_{\varepsilon \downarrow 0} \varepsilon N(T, \varepsilon) = L^S(T) \); whereas the strong law of large numbers gives

\[
\lim_{\varepsilon \downarrow 0} \frac{1}{N(T, \varepsilon)} \sum_{j=1}^{N(T, \varepsilon)} \xi_{\ell_j} = \mathbb{E}(\xi_1).
\]

Back into (3.17) and with the help of (3.18), these considerations give

\[
Z(T) S(T) = \int_0^T Z(t) \, dS(t) + \mathbb{E}(\xi_1) \cdot L^S(T), \quad 0 \leq T < \infty,
\]

that is, (3.5).

\[
\text{4 Conventional Reflection}
\]

In a similar manner one can establish the following analogue of Theorem 1.11, which uses the conventional reflection in place of the Skorokhod reflection.

**Theorem 4.1.** Fix a constant \( \alpha \in (0, 1) \). There exists an enlargement \((\hat{\Omega}, \hat{\mathcal{F}}, \hat{\mathbb{P}}), \hat{\mathcal{F}} = \{\hat{F}(t)\}_{0 \leq t < \infty}\) of the filtered probability space \((\Omega, \mathcal{F}, \mathbb{P}), \mathcal{F} = \{\mathcal{F}(t)\}_{0 \leq t < \infty}\) with a measure-preserving map \( \pi : \Omega \to \hat{\Omega} \), and on this enlarged space a continuous semimartingale \( \hat{X}(\cdot) \) that satisfies

\[
|\hat{X}(\cdot)| = |U(\cdot)|, \quad L\hat{X}(\cdot) = \alpha L|U|(\cdot), \quad \hat{X}(\cdot) = \int_0^\cdot \mathbb{sgn}(\hat{X}(t)) \, d\hat{U}(t) + \frac{2\alpha - 1}{\alpha} L\hat{X}(\cdot).
\]

Here

\[
\hat{U}(\cdot) := \int_0^\cdot \mathbb{sgn}(U(t)) \, dU(t)
\]

(4.2)

is the Lévy transform of the semimartingale \( U(\cdot) \), and the classical reflection \( R(\cdot) = |U(\cdot)| \) of \( U(\cdot) \) coincides with the Skorokhod reflection of the process \( \hat{U}(\cdot) \) in (4.2), namely

\[
\hat{S}(t) := \hat{U}(t) + \max_{0 \leq s \leq t} \left( -\hat{U}(s) \right), \quad 0 \leq t < \infty.
\]
Remark and

4.1. It is well-known (see, for instance, C

with parameter

the S

as the local time property

(4.2). The property

(4.1) follows from this last display. On the other hand, since the local time

here by

\( \hat{\tau} \), \( \hat{\tau} \) becomes

without any additional conditions on \( U(\cdot) \). Thus, the stochastic integral equation of (4.4)

becomes

\[
\hat{\tau}(\cdot) = \int_0^\cdot \mathrm{sgn}(\hat{\tau}(t)) \, dU(t) + 2\gamma - 1 \gamma \, \hat{\tau}(\cdot),
\]

and observed that it is another standard Brownian motion. Thus, the stochastic integral equation of (4.4)

becomes

\[
\hat{\tau}(\cdot) = \int_0^\cdot \mathrm{sgn}(\hat{\tau}(t)) \, dW(t) + 2\gamma - 1 \gamma \, \hat{\tau}(\cdot) = \hat{\tau}(\cdot) + 2\gamma - 1 \gamma \, \hat{\tau}(\cdot)
\]

with \( \hat{\tau}(\cdot) = \int_0^\cdot \mathrm{sgn}(\hat{\tau}(t)) \, dW(t) \) yet another standard Brownian motion.

The HARRISON-SHEPP (1981) theory characterizes now \( \hat{\tau}(\cdot) \) as skew Brownian motion with skewness parameter \( \alpha \). The processes \( \hat{\tau}(\cdot) \) and \( \hat{\tau}(\cdot) \) generate the same filtration, as do the processes

\[
\hat{\tau}(\cdot) = \int_0^\cdot \mathrm{sgn}(\hat{\tau}(t)) \, d\hat{\tau}(t) = W(\cdot)
\]

and the first filtration is finer than the second.
4.1 Skew Bessel Processes

In this subsection suppose that \( U^2(\cdot) \) is a squared Bessel process with dimension \( \delta \in (1, 2) \), i.e., \( U^2(\cdot) \) is the unique strong solution of the equation

\[
U^2(t) = \delta t + 2 \int_0^t \sqrt{U^2(s)} \, dB(s), \quad 0 \leq t < \infty
\]

for some standard, real-valued Brownian motion \( B(\cdot) \). When \( \delta \in (1, 2) \), the square root \( R(\cdot) := |U(\cdot)| \geq 0 \) of this process is a semimartingale that keeps visiting the origin almost surely, and can be decomposed as

\[
R(\cdot) = \int_0^\cdot \frac{\delta - 1}{2} \frac{1}{R(t)} \, 1_{\{R(t) \neq 0\}} \, dt + B(\cdot) \quad \text{with} \quad L^R(\cdot) \equiv 0, \quad \int_0^\cdot 1_{\{R(t) = 0\}} \, dt \equiv 0. \quad (4.5)
\]

For the study of the stochastic differential equation (4.5) with \( \delta \in (1, 2) \) see, for example, Cherny (2000).

Given \( \alpha \in (0, 1) \), following again the argument of the proof in section 3 through verbatim, with \( S(\cdot), X(\cdot) \) replaced respectively by \( R(\cdot), \tilde{X}(\cdot) \), we unfold the nonnegative Bessel process \( R(\cdot) \) to obtain

\[
\tilde{X}(\cdot) = Z(\cdot)R(\cdot) = \int_0^\cdot Z(t) \, dR(t) + (2\alpha - 1)L^R(\cdot) = \int_0^\cdot \frac{\delta - 1}{2} \frac{1}{\tilde{X}(t)} \, 1_{\{\tilde{X}(t) \neq 0\}} \, dt + \tilde{\beta}(\cdot), \quad (4.6)
\]

with \( Z(\cdot) = \text{sgn}(\tilde{X}(\cdot)) \) and with \( \tilde{\beta}(\cdot) := \int_0^\cdot Z(t) \, dB(t) \) another standard Brownian motion on an extended probability space, as a consequence of Theorem 4.1 and of the properties in (4.5). We note that the semimartingale \( \tilde{X}(\cdot) \) does not accumulate local time at the origin, because of \( L^R(\cdot) \equiv 0 \).

We claim that the process \( \tilde{X}(\cdot) \) constructed here in (4.6) is the \( \delta \)-dimensional skew Bessel process with skewness parameter \( \alpha \). This process was introduced and studied in Blei (2012).

Indeed, let us consider the functions \( g(x) := |x|^{2-\delta}/(2-\delta) \) and \( G(x) := \text{sgn}(x) \cdot g(x) \) for \( x \in \mathbb{R} \), and examine \( g(\tilde{X}(\cdot)) \) and \( G(\tilde{X}(\cdot)) \). This scaling is a right choice to measure the boundary behavior of \( \tilde{X}(\cdot) \) around the origin. By substituting \( q = 2 - \delta, \quad p = (2 - \delta)/(1 - \delta), \quad \nu = -1/2 \) in Proposition XI.11 of Revuz & Yor (2005), we find there exists a (nonnegative) one-dimensional Bessel process \( \rho(\cdot) \) on the same probability space such that \( \rho(0) = (2 - \delta)^{1-\delta} \, g(\tilde{X}(0)) \) and

\[
g(\tilde{X}(t)) = \frac{1}{2 - \delta} |\tilde{X}(t)|^{2-\delta} = \frac{1}{(2 - \delta)^{\delta-1}} \rho(\Lambda(t)), \quad 0 \leq t < \infty,
\]

where

\[
\Lambda(t) := \inf\{s \geq 0 : K(s) \geq t\}, \quad K(s) := \int_0^s \left(\rho(u)\right)^{\frac{3(\delta-2)}{2}} \, du,
\]

that is, \( g(\tilde{X}(\cdot)) \) is a time-changed, conventionally reflected Brownian motion with the stochastic clock \( \Lambda(\cdot) \). Thus the local time of \( g(\tilde{X}(\cdot)) \) accumulates at the origin with this clock \( \Lambda(\cdot) \).

In the same manner as in the construction of \( Z(\cdot)R(\cdot) \) in Theorem 4.1, we obtain here

\[
G(\tilde{X}(T)) = \text{sgn}(\tilde{X}(T))g(\tilde{X}(T)) = \int_0^T \text{sgn}(\tilde{X}(t)) \, d(\rho(\tilde{X}(t))) + (2\alpha - 1) L^G(\tilde{X})(T)
\]

as well as

\[
L^G(\tilde{X})(\cdot) - L^{-G}(\tilde{X})(\cdot) = (2\alpha - 1) (L^G(\tilde{X})(\cdot) + L^{-G}(\tilde{X})(\cdot)) \quad (4.7)
\]

and

\[
(1 - \alpha) L^{G(\tilde{X})}(\cdot) = \alpha L^{-G(\tilde{X})}(\cdot), \quad L^G(\tilde{X})(\cdot) = \frac{1}{2} (L^G(\tilde{X})(\cdot) + L^{-G(\tilde{X})}(\cdot)).
\]
in the notation of (1.5). From these relationships (4.7), and on the strength of Theorem 2.22 of BLEI (2012),
we identify the process of (4.6) as the \(\delta\)-dimensional skew BESSEL process. Here the process \(G(\tilde{X}(\cdot))\)
and its local time \(L_G(\tilde{X}(\cdot))\) correspond to \(Y(\cdot)\) and \(L^\alpha(\cdot)\), respectively, in the notation of BLEI (2012).

For various properties and representations of this process, we refer the study of BLEI (2012), in particular,
Remark 2.26 there.

5 An Application: Two Diffusive Particles with Asymmetric Collisions

In the paper [FIKP], the authors construct a planar continuous semimartingale \(\mathcal{X}(\cdot) = (X_1(\cdot), X_2(\cdot))\) with
dynamics
\[
\begin{align*}
dX_1(t) &= \left(gX_1(t)I_{X_1(t) \leq X_2(t)} - h \delta X_1(t) \mathbb{1}_{X_1(t) > X_2(t)}\right)dt + \left(\rho X_1(t) \mathbb{1}_{X_1(t) > X_2(t)} + \sigma \delta X_1(t) \mathbb{1}_{X_1(t) \leq X_2(t)}\right)dB_1(t), \\
dX_2(t) &= \left(gX_2(t)I_{X_1(t) > X_2(t)} - h \delta X_2(t) \mathbb{1}_{X_1(t) \leq X_2(t)}\right)dt + \left(\rho X_2(t) \mathbb{1}_{X_1(t) > X_2(t)} + \sigma \delta X_2(t) \mathbb{1}_{X_1(t) \leq X_2(t)}\right)dB_2(t),
\end{align*}
\]
for arbitrary real constants \(g, h, \rho > 0, \sigma > 0\) with \(\rho^2 + \sigma^2 = 1\). They show that, for an arbitrary
initial condition \((X_1(0), X_2(0)) = (x_1, x_2) \in \mathbb{R}^2\) and with \((B_1(\cdot), B_2(\cdot))\) a planar Brownian motion, the
system of (5.1) has a pathwise solution.

This is a model for two “competing” Brownian particles, with diffusive motions whose drift and dispersion
characteristics are assigned according to their ranks.

- In another recent paper FERNHOLZ, ICHIBA & KARATZAS (2013), hereafter referred to as [FIK], a planar
continuous semimartingale \(\tilde{\mathcal{X}}(\cdot) = (\tilde{X}_1(\cdot), \tilde{X}_2(\cdot))\) is constructed according to the dynamics
\[
\begin{align*}
d\tilde{X}_1(t) &= \left(gX_1(t)I_{\tilde{X}_1(t) \leq \tilde{X}_2(t)} - h \delta \tilde{X}_1(t) \mathbb{1}_{\tilde{X}_1(t) > \tilde{X}_2(t)}\right)dt + \left(\rho X_1(t) \mathbb{1}_{\tilde{X}_1(t) > \tilde{X}_2(t)} + \sigma \delta \tilde{X}_1(t) \mathbb{1}_{\tilde{X}_1(t) \leq \tilde{X}_2(t)}\right)d\tilde{B}_1(t) \\
&\quad + \frac{1 - \zeta_1}{2} dL_{\tilde{X}_1, \tilde{X}_2}(t), \\
d\tilde{X}_2(t) &= \left(gX_2(t)I_{\tilde{X}_1(t) > \tilde{X}_2(t)} - h \delta \tilde{X}_2(t) \mathbb{1}_{\tilde{X}_1(t) \leq \tilde{X}_2(t)}\right)dt + \left(\rho X_2(t) \mathbb{1}_{\tilde{X}_1(t) > \tilde{X}_2(t)} + \sigma \delta \tilde{X}_2(t) \mathbb{1}_{\tilde{X}_1(t) \leq \tilde{X}_2(t)}\right)d\tilde{B}_2(t) \\
&\quad + \frac{1 - \zeta_2}{2} dL_{\tilde{X}_1, \tilde{X}_2}(t),
\end{align*}
\]
Here again \(g, h\) are arbitrary real constants, \(\rho > 0\) and \(\sigma > 0\) satisfy \(\rho^2 + \sigma^2 = 1\), whereas \(\zeta_1, \eta_1\) are
real constants satisfying
\[
0 \leq \alpha := \frac{\eta}{\eta + \zeta} \leq 1, \quad \zeta := 1 + \frac{\zeta_1 - \zeta_2}{2}, \quad \eta := 1 - \frac{\eta_1 - \eta_2}{2}, \quad \zeta + \eta \neq 0.
\]
This new system is a version of the previous competing Brownian particle system, but now with elastic and asymmetric collisions whose effect is modeled by the local time terms \(L_{\tilde{X}_2, \tilde{X}_1}(\cdot)\) and \(L_{\tilde{X}_2, \tilde{X}_1}(\cdot)\). Every time the two particles collide, their trajectories feel a “drag” proportional to these local time terms, whose
presence makes the analysis of the system (5.3) considerably more involved than that of (5.1).

It is shown in [FIK] under the above conditions that, for an arbitrary initial condition \((\tilde{X}_1(0), \tilde{X}_2(0)) = (x_1, x_2) \in \mathbb{R}^2\), and with \((\tilde{B}_1(\cdot), \tilde{B}_2(\cdot))\) a planar Brownian motion, the system of (5.3) has a pathwise
unique, strong solution.

- We shall show how to use the unfolding of Theorem 1.1 in order to construct the planar process \(\tilde{\mathcal{X}}(\cdot) = (\tilde{X}_1(\cdot), \tilde{X}_2(\cdot))\) of (5.3) with skew-elastic collisions, starting from the planar diffusion \(\mathcal{X}(\cdot) = (X_1(\cdot), X_2(\cdot))\) of (5.1). For simplicity, we shall take the initial condition \((x_1, x_2) = (0, 0)\) from now on.
Theorem 5.1. Suppose we are given a planar continuous semimartingale \( X(\cdot) = (X_1(\cdot), X_2(\cdot)) \) that satisfies the system of (5.1), (5.2) on some filtered probability space \( (\Omega, \mathcal{F}, \mathbb{P}) \), \( \mathbb{F} = \{ \mathcal{F}(t) \}_{0 \leq t < \infty} \) with a planar Brownian motion \( (B_1(\cdot), B_2(\cdot)) \).

There exists then an enlargement \( (\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}}) = \{ \tilde{\mathcal{F}}(t) \}_{0 \leq t < \infty} \) of this filtered probability space, with a planar Brownian motion \( (\tilde{B}_1(\cdot), \tilde{B}_2(\cdot)) \), and on it a planar continuous semimartingale \( \tilde{X}(\cdot) = (\tilde{X}_1(\cdot), \tilde{X}_2(\cdot)) \) that satisfies the system of (5.3), (5.4) with skew-elastic collisions, as well as

\[
(X_1(t) - X_2(t)) + \sup_{0 \leq s \leq t} (X_1(s) - X_2(s))^+ = |\tilde{X}_1(t) - \tilde{X}_2(t)|, \quad 0 \leq t < \infty.
\]

In other words, the size of the gap between the new processes \( \tilde{X}_1(\cdot), \tilde{X}_2(\cdot) \) coincides with the Skorokhod reflection of the difference \( X_1(\cdot) - X_2(\cdot) \) of the original processes about the origin. We devote the remainder of this section to the proof of this result.

5.1 Reduction to symmetric local times

First, some preparatory steps. We define the averages \( \tilde{\zeta} := (\zeta_1 + \zeta_2)/2 \), \( \tilde{\eta} := (\eta_1 + \eta_2)/2 \), and introduce yet another parameter

\[
\beta := \alpha \cdot \frac{\zeta_1 + \zeta_2}{2} + (1 - \alpha) \cdot \frac{\eta_1 + \eta_2}{2} = \alpha \tilde{\zeta} + (1 - \alpha) \tilde{\eta}.
\]

For notational simplicity we shall write all the processes related to the skew collisions with a tilde, e.g., \( \tilde{Y}(\cdot) := \tilde{X}_1(\cdot) - \tilde{X}_2(\cdot) \). From the relation between the right local time \( L^\tilde{Y}(\cdot) \) and the symmetric local time \( \hat{L}^\tilde{Y}(\cdot) \) as in (1.5), we obtain the relations

\[
\zeta L^\tilde{Y}(\cdot) = \eta L^{-\tilde{Y}}(\cdot), \quad L^\tilde{Y}(\cdot) = 2\alpha \hat{L}^\tilde{Y}(\cdot), \quad L^{-\tilde{Y}}(\cdot) := L^{-\tilde{Y}}(\cdot) = 2(1 - \alpha) \hat{L}^\tilde{Y}(\cdot)
\]

as in [FK]. This way, the system (5.3)-(5.4) can be recast as

\[
d\tilde{X}_1(t) = (g1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}} - h1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}})dt + (\rho1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}} + \sigma1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}})d\tilde{B}_1(t)
\]
\[
+ (2\alpha - \beta) d\hat{L}^\tilde{Y}(t),
\]

\[
d\tilde{X}_2(t) = (g1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}} - h1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}})dt + (\rho1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}} + \sigma1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}})d\tilde{B}_2(t)
\]
\[
+ (2 - 2\alpha - \beta) d\hat{L}^\tilde{Y}(t).
\]

We shall construct the system (5.7)-(5.8) first, and then obtain from it the system (5.3)-(5.4).

5.2 Proof of Theorem 5.1

By applying a Girsanov change of measure twice, we can remove the drifts from both of the systems (5.1)-(5.2) and (5.7)-(5.8). Then, in the following, let us construct the two-dimensional Brownian motion with rank-based dispersions and skew-elastic collisions

\[
d\tilde{X}_1(t) = (\rho1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}} + \sigma1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}})d\tilde{B}_1(t) + (2\alpha - \beta) d\hat{L}^\tilde{Y}(t),
\]

\[
d\tilde{X}_2(t) = (\rho1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}} + \sigma1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}})d\tilde{B}_2(t) + (2 - 2\alpha - \beta) d\hat{L}^\tilde{Y}(t)
\]
from the solution \( ((X_1(\cdot), X_2(\cdot)), (B_1(\cdot), B_2(\cdot))) \) of the system

\[
\begin{align*}
    dX_1(t) &= (\rho 1_{\{X_1(t)>0\}} + \sigma 1_{\{X_1(t)\leq 0\}}) dB_1(t), \\
    dX_2(t) &= (\rho 1_{\{X_1(t)\leq 0\}} + \sigma 1_{\{X_1(t)>0\}}) dB_2(t),
\end{align*}
\]

which is known from [FIKP] to be strongly solvable. Since there is no drift in these last equations, the difference \( Y(\cdot) := X_1(\cdot) - X_2(\cdot) \) between the two components of the system (5.10) is given by the real-valued Brownian motion

\[
Y(\cdot) = W(\cdot) := \rho W_1(\cdot) + \sigma W_2(\cdot).
\]

Here

\[
\begin{align*}
    W_1(\cdot) &= \int_0^\cdot 1_{\{X_1(t)>X_2(t)\}} dB_1(t) - \int_0^\cdot 1_{\{X_1(t)\leq X_2(t)\}} dB_2(t), \\
    W_2(t) &= \int_0^\cdot 1_{\{X_1(t)\leq X_2(t)\}} dB_1(t) - \int_0^\cdot 1_{\{X_1(t)>X_2(t)\}} dB_2(t)
\end{align*}
\]

are independent Brownian motions. As in [FIKP], let us recall also the Brownian motion

\[
V(\cdot) := \rho V_1(\cdot) + \sigma V_2(\cdot),
\]

where again

\[
\begin{align*}
    V_1(\cdot) &= \int_0^\cdot 1_{\{X_1(t)>X_2(t)\}} dB_1(t) + \int_0^\cdot 1_{\{X_1(t)\leq X_2(t)\}} dB_2(t), \\
    V_2(\cdot) &= \int_0^\cdot 1_{\{X_1(t)\leq X_2(t)\}} dB_1(t) + \int_0^\cdot 1_{\{X_1(t)>X_2(t)\}} dB_2(t)
\end{align*}
\]

are independent Brownian motions. For a given number \( \alpha \in (0, 1) \), there exists by Theorem 1.1 an adapted, continuous process \( \tilde{Y}(\cdot) \) which satisfies

\[
Y(t) + \sup_{0 \leq s \leq t} (Y(s))^+ = |\tilde{Y}(t)|, \quad 0 \leq t < \infty
\]

as well as

\[
\tilde{Y}(\cdot) = \int_0^\cdot \text{sgn}(\tilde{Y}(t)) \, dY(t) + \frac{2\alpha - 1}{\alpha} L\tilde{Y}(\cdot) = \int_0^\cdot \text{sgn}(\tilde{Y}(t)) \, dW(t) + 2(2\alpha - 1) \tilde{\tilde{L}}\tilde{Y}(\cdot),
\]

where the last equality follows from Proposition 2.1 and 5.6. Thus, the “unfolded process” \( \tilde{Y}(\cdot) \) is a skew Brownian motion, with skewness parameter \( \alpha \).

Now let us define the new planar Brownian motion \( (\tilde{B}_1(\cdot), \tilde{B}_2(\cdot)) \) as

\[
\begin{align*}
    d\tilde{B}_1(\cdot) &= (1_{\{Y(\cdot)>0, \tilde{Y}(\cdot)>0\}} - 1_{\{Y(\cdot)=0, \tilde{Y}(\cdot)>0\}}) dB_1(\cdot) + (1_{\{Y(\cdot)>0, \tilde{Y}(\cdot)=0\}} - 1_{\{Y(\cdot)=0, \tilde{Y}(\cdot)>0\}}) dB_2(\cdot), \\
    d\tilde{B}_2(\cdot) &= (1_{\{Y(\cdot)>0, \tilde{Y}(\cdot)=0\}} - 1_{\{Y(\cdot)=0, \tilde{Y}(\cdot)>0\}}) dB_1(\cdot) + (1_{\{Y(\cdot)>0, \tilde{Y}(\cdot)>0\}} - 1_{\{Y(\cdot)=0, \tilde{Y}(\cdot)=0\}}) dB_2(\cdot),
\end{align*}
\]

and, with the number \( \beta \in \mathbb{R} \) as in (5.5), the processes \( \tilde{\Xi}(\cdot), (\tilde{X}_1(\cdot), \tilde{X}_2(\cdot)) \) and \( (\tilde{V}(\cdot), \tilde{W}(\cdot)) \) by

\[
\begin{align*}
    \tilde{\Xi}(\cdot) &= \tilde{V}(\cdot) + 2(1-\beta)\tilde{\tilde{L}}\tilde{Y}(\cdot), \quad \tilde{X}_1(\cdot) := \frac{\tilde{\Xi}(\cdot) + \tilde{Y}(\cdot)}{2}, \quad \tilde{X}_2(\cdot) := \frac{\tilde{\Xi}(\cdot) - \tilde{Y}(\cdot)}{2}, \\
    d\tilde{V}(\cdot) &= (\rho 1_{\{\tilde{Y}(\cdot)>0\}} + \sigma 1_{\{\tilde{Y}(\cdot)=0\}}) d\tilde{B}_1(\cdot) + (\rho 1_{\{\tilde{Y}(\cdot)=0\}} + \sigma 1_{\{\tilde{Y}(\cdot)>0\}}) d\tilde{B}_2(\cdot), \\
    d\tilde{W}(\cdot) &= (\rho 1_{\{\tilde{Y}(\cdot)>0\}} + \sigma 1_{\{\tilde{Y}(\cdot)=0\}}) d\tilde{B}_1(\cdot) - (\rho 1_{\{\tilde{Y}(\cdot)=0\}} + \sigma 1_{\{\tilde{Y}(\cdot)>0\}}) d\tilde{B}_2(\cdot),
\end{align*}
\]
Then by (5.11) and (5.13) we obtain

\[
\text{sgn}(\tilde{Y}(\cdot))dW(\cdot) = \text{sgn}(\tilde{Y}(\cdot)) \left[ (\rho 1_{\{Y(\cdot) > X(\cdot)\}} + \sigma 1_{\{X(\cdot) \leq X(\cdot)\}})dB_1(\cdot) - (\rho 1_{\{Y(\cdot) \leq X(\cdot)\}} + \sigma 1_{\{X(\cdot) > X(\cdot)\}})dB_2(\cdot) \right]
\]

\[
= \text{sgn}(\tilde{Y}(\cdot)) \left[ (\rho 1_{\{Y(\cdot) > 0\}} + \sigma 1_{\{X(\cdot) \leq 0\}})dB_1(\cdot) - (\rho 1_{\{Y(\cdot) \leq 0\}} + \sigma 1_{\{X(\cdot) > 0\}})dB_2(\cdot) \right] ,
\]

\[
d\tilde{W}(\cdot) = (\rho 1_{\{\tilde{Y}(\cdot) > 0\}} + \sigma 1_{\{\tilde{Y}(\cdot) \leq 0\}})dB_1(\cdot) - (\rho 1_{\{\tilde{Y}(\cdot) \leq 0\}} + \sigma 1_{\{\tilde{Y}(\cdot) > 0\}})dB_2(\cdot)
\]

\[
= 1_{\{\tilde{Y}(\cdot) > 0\}} (\rho dB_1(\cdot) - \sigma dB_2(\cdot)) + 1_{\{\tilde{Y}(\cdot) \leq 0\}} (\sigma dB_1(\cdot) - \rho dB_2(\cdot)) .
\]

Because of the relationship between \((B_1(\cdot), B_2(\cdot))\) and \((\tilde{B}_1(\cdot), \tilde{B}_2(\cdot))\), it can be shown that

\[
d\tilde{W}(\cdot) = \text{sgn}(\tilde{Y}(\cdot)) dW(\cdot) . \tag{5.14}
\]

In fact, these identities can be verified formally via the following table:

| signs of \((Y(\cdot), \tilde{Y}(\cdot))\) | \(dB_1(\cdot)\) | \(dB_2(\cdot)\) | \(dW(\cdot)\) | \(= \text{sgn}(Y(\cdot))dW(\cdot)\) |
|---|---|---|---|---|
| (+, +) | \(dB_1(\cdot)\) | \(dB_2(\cdot)\) | \(\rho dB_1(\cdot) - \sigma dB_2(\cdot)\) | \(\rho dB_1(\cdot) - \sigma dB_2(\cdot)\) |
| (−, +) | −\(dB_2(\cdot)\) | −\(dB_1(\cdot)\) | \(\rho dB_1(\cdot) - \sigma dB_2(\cdot)\) | \(\sigma dB_1(\cdot) - \rho dB_2(\cdot)\) |
| (+, −) | \(dB_2(\cdot)\) | \(dB_1(\cdot)\) | \(\sigma dB_1(\cdot) - \rho dB_2(\cdot)\) | \(\rho dB_1(\cdot) + \sigma dB_2(\cdot)\) |
| (−, −) | −\(dB_1(\cdot)\) | −\(dB_2(\cdot)\) | \(\sigma dB_1(\cdot) - \rho dB_2(\cdot)\) | \(\rho dB_1(\cdot) + \sigma dB_2(\cdot)\) |

Substituting this relation (5.14) into (5.12) and recalling (5.13), we obtain

\[
d(\tilde{X}_1(t) - \tilde{X}_2(t)) = d\tilde{Y}(t) = d\tilde{W}(t) + 2(2\alpha - 1)d\tilde{L}\tilde{Y}(t) . \tag{5.15}
\]

Moreover, because of the correspondence between \((\tilde{V}(\cdot), \tilde{W}(\cdot))\) and \((V(\cdot), W(\cdot))\) and the relation (5.13), we obtain

\[
\frac{1}{2} d(\tilde{V}(t) + \tilde{W}(t)) = (\rho 1_{\{\tilde{Y}(\cdot) > 0\}} + \sigma 1_{\{\tilde{Y}(\cdot) \leq 0\}}) dB_1(t) , \tag{5.16}
\]

\[
\frac{1}{2} d(\tilde{V}(t) - \tilde{W}(t)) = (\sigma 1_{\{\tilde{Y}(\cdot) > 0\}} + \rho 1_{\{\tilde{Y}(\cdot) \leq 0\}}) dB_2(t) .
\]

Therefore, by calculating the coefficients in front of the local time terms and by combining (5.13), (5.15) and (5.16), we can verify that \((\tilde{X}_1(\cdot), \tilde{X}_2(\cdot))\) satisfies

\[
d\tilde{X}_1(t) = (\rho 1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}} + \sigma 1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}}) dB_1(t) + (2\alpha - \beta)d\tilde{L}\tilde{Y}(t) , \tag{5.17}
\]

\[
d\tilde{X}_2(t) = (\rho 1_{\{\tilde{X}_1(t) \leq \tilde{X}_2(t)\}} + \sigma 1_{\{\tilde{X}_1(t) > \tilde{X}_2(t)\}}) dB_2(t) + (2 - 2\alpha - \beta)d\tilde{L}\tilde{Y}(t)
\]

that is, (5.9) with the new Brownian motion \((\tilde{B}_1(\cdot), \tilde{B}_2(\cdot))\).

By the Girsanov theorem, we obtain (5.7)–(5.8); whereas the relationship (5.6) between the left local time \(L^{-}\tilde{Y}(\cdot)\) and the right local time \(L^{+}\tilde{Y}(\cdot)\) allows us now to recover the dynamics of (5.3)–(5.4) from those of (5.1)–(5.2).
6 Appendix: Proof of Proposition 2.3

Given a planar Brownian motion \((B_1(\cdot), B_2(\cdot))\) on a probability space \((\Omega, \mathcal{F}, \mathbb{P})\) and real constants \(\alpha \in (0, 1)\), \(x_0 \in \mathbb{R}\), we shall construct a process \(X(\cdot) := q(Y(\cdot))\) from the solution \(Y(\cdot)\) of the stochastic differential equation

\[ Y(\cdot) = p(x_0) + \int_0^\cdot s(Y(t)) \, dB_1(t) + B_2(t), \]

where \(p(\cdot), q(\cdot)\) and \(s(\cdot)\) are defined by

\[ p(x) := (1 - \alpha)x \mathbf{1}_{(0, \infty)}(x) + \alpha x \mathbf{1}_{(-\infty, 0]}(x), \quad q(x) := \frac{1}{1 - \alpha} \mathbf{1}_{(0, \infty)}(x) + \frac{1}{\alpha} \mathbf{1}_{(-\infty, 0]}(x), \]

\[ s(x) := (1 - \alpha) \mathbf{1}_{(0, \infty)}(x) + \alpha \mathbf{1}_{(-\infty, 0]}(x); \quad x \in \mathbb{R}. \]

From the work on ANAKA (1972) we know that the equation (6.1) has a pathwise unique, strong solution. Since \(q(p(x)) = x\), \(x \in \mathbb{R}\), by applying the Itô-TANAKA formula to the process \(X(\cdot) = q(Y(\cdot))\) we identify the dynamics of \(X(\cdot)\) as those of the skew Brownian motion (HARRISON & SHEPP (1981)), namely

\[ X(\cdot) = x_0 + (B_1(\cdot) + B_2(\cdot)) + \frac{2\alpha - 1}{\alpha} L^X(\cdot), \]

driven by the Brownian motion \((B_1(\cdot) + B_2(\cdot))\). We rewrite this equation in the form

\[ X(\cdot) - x_0 - \int_0^\cdot \text{sgn}(X(t)) \, dU(t) - V(\cdot) = \frac{2\alpha - 1}{\alpha} L^X(\cdot) = 2(2\alpha - 1) \tilde{L}^X(\cdot) \]

of (2.16), driven by a new planar Brownian motion \((U(\cdot), V(\cdot))\) with components

\[ U(\cdot) := \int_0^\cdot \mathbf{1}_{\{X(t) > 0\}} \, dB_1(t) - \int_0^\cdot \mathbf{1}_{\{X(t) \leq 0\}} \, dB_2(t), \]

\[ V(\cdot) := \int_0^\cdot \mathbf{1}_{\{X(t) \leq 0\}} \, dB_1(t) + \int_0^t \mathbf{1}_{\{X(t) > 0\}} \, dB_2(t). \]

Therefore, the perturbed skew TANAKA equation (2.16) has the weak solution \((X(\cdot), (U(\cdot), V(\cdot)))\) just constructed.

Conversely, suppose we start with an arbitrary weak solution \((X(\cdot), (U(\cdot), V(\cdot)))\) of the equation (2.16), with \((U(\cdot), V(\cdot))\) a planar Brownian motion. Then we can cast this equation in the form (6.2) in terms of the planar Brownian motion \((B_1(\cdot), B_2(\cdot))\) whose components are given by "disentangling" in (6.3), (6.4), namely

\[ B_1(\cdot) = \int_0^\cdot \mathbf{1}_{\{X(t) > 0\}} dU(t) + \int_0^\cdot \mathbf{1}_{\{X(t) \leq 0\}} dV(t), \]

\[ B_2(\cdot) = \int_0^\cdot \mathbf{1}_{\{X(t) > 0\}} dV(t) - \int_0^\cdot \mathbf{1}_{\{X(t) \leq 0\}} dU(t). \]

But this shows that \(X(\cdot)\) is skew Brownian motion, so its probability distribution is determined uniquely.

In other words, the equation of (2.16) admits a weak solution, and this solution in unique in the sense of the probability distribution.

- Now we shall see that we have not just uniqueness in distribution, but also pathwise uniqueness, for the equation (2.16) driven by the planar Brownian motion \((U(\cdot), V(\cdot))\). The argument that follows is based on Lemma 1 of LE GALL (1983), and is almost identical to the proof of Theorem 8.1 of [FIKP] except

...
for the evaluation of the additional local times. Note that Le Gall’s Lemma 1 works for continuous semimartingales, in general.

Suppose that there are two solutions $X_1(\cdot)$ and $X_2(\cdot)$ of (2.16), defined on the same probability space as the driving planar Brownian motion $(U(\cdot), V(\cdot))$. We shall check their difference $D(\cdot) := X_1(\cdot) - X_2(\cdot)$ satisfies (c.f. (8.4) in Fernholz et al. (2011)):

$$\mathbb{E} \left[ \int_0^T \frac{d\langle D \rangle(s)}{D(s)} 1_{\{D(s) > 0\}} \right] < \infty, \quad 0 < T < \infty,$$

where

$$\langle D \rangle(\cdot) = \int_0^\infty \left( \text{sgn}(X_1(t)) - \text{sgn}(X_2(t)) \right)^2 dt \leq 2 \int_0^\infty |\text{sgn}(X_1(t)) - \text{sgn}(X_2(t))| dt.$$

We approximate the signum function by a sequence $\{f_k\}_{k \in \mathbb{N}} \subset C^1(\mathbb{R})$ which converges to the function $f_\infty(\cdot) = \text{sgn}(\cdot)$ pointwise and satisfies $\lim_{k \to \infty} \|f_k\|_{TV} = \|f_\infty\|_{TV}$. Now the parametrized process

$$Z^{(u)}(t) := (1-u)X_1(t) + uX_2(t), \quad 0 \leq u \leq 1, \quad 0 \leq t < \infty$$

takes the form of

$$Z^{(u)}(\cdot) = x_0 + \int_0^\cdot \left( (1-u)\text{sgn}(X_1(t)) + u\text{sgn}(X_2(t)) \right) dU(t) + V(\cdot) + \frac{2\alpha - 1}{\alpha} (uL_{X_1}(\cdot) + (1-u)L_{X_2}(\cdot)).$$

The local times in the last term do not affect the size of $\langle Z^{(u)}(\cdot) \rangle$, for which we have the estimate $\mathbb{E}(\langle Z^{(u)}(T) \rangle) \leq 2T$. Proceeding as in [FIKP] we obtain for every $\delta > 0$ the bound

$$\mathbb{E} \left[ \int_0^T \frac{|f_k(X_1(s)) - f_k(X_2(s))|}{X_1(s) - X_2(s)} 1_{\{X_1(s) - X_2(s) > \delta\}} ds \right] \leq c \|f_k\|_{TV} \sup_{a,u} \mathbb{E} (2L^{(u)}(T, a)),$$

where $L^{(u)}(T, a)$ is the right local time of the continuous semimartingale $Z^{(u)}(\cdot)$ accumulated at $a \in \mathbb{R}$ and $c$ is a constant chosen independently of $k, u, \delta$. Letting $k \uparrow \infty$ and $\delta \downarrow 0$, we estimate

$$\mathbb{E} \left[ \int_0^T \frac{d\langle D \rangle(s)}{D(s)} 1_{\{D(s) > 0\}} \right] < 2c \|f_\infty\|_{TV} \sup_{a,u} \mathbb{E} (2L^{(u)}(T, a)).$$

Finally, we estimate $\mathbb{E}(L^{(u)}(T, a))$ using Tanaka’s formula

$$|Z^{(u)}(T) - a| = |Z^{(u)}(0) - a| + \int_0^T \text{sgn}(Z^{(u)}(t) - a) dZ^{(u)}(t) + 2L^{(u)}(T, a),$$

and a combination of the Cauchy-Schwartz inequality and the Itô’s isometry:

$$\mathbb{E}(2L^{(u)}(T, a)) \leq \mathbb{E}|Z^{(u)}(T) - Z^{(u)}(0)| + \{\mathbb{E}(\langle Z^{(u)}(T) \rangle)\}^{1/2}$$

$$+ \frac{2\alpha - 1}{\alpha} \left( u\mathbb{E}(L_{X_1}(T)) + (1-u)\mathbb{E}(L_{X_2}(T)) \right)$$

$$\leq 2 \left[ \{\mathbb{E}(\langle Z^{(u)}(T) \rangle)\}^{1/2} + \frac{2\alpha - 1}{\alpha} \left( u\mathbb{E}(L_{X_1}(T)) + (1-u)\mathbb{E}(L_{X_2}(T)) \right) \right].$$
The last term \( \mathbb{E}(L^X_i(T)) \) is evaluated by the same procedure: by TANAKA’s formula
\[
\frac{1}{\alpha}L^X_i(T) = |X_i(T)| - |X_1(0)| - \int_0^T \text{sgn}(X_i(t)) \, dV(t) - U(T),
\]
and hence
\[
\mathbb{E}(L^X_i(T)) \leq 2\alpha \{\mathbb{E}(\langle X_i \rangle(T))\}^{1/2} \leq 2^{3/2} \alpha T^{1/2}, \quad i = 1, 2.
\]
Therefore, we obtain (6.5), and by Lemma 1 of LE Gall (1983) we verify \( L^D(\cdot) = L^{X_1-X_2}(\cdot) \equiv 0 \).

\textbf{Final step:} By exchanging the rôles of \( X_1(\cdot) \) and \( X_2(\cdot) \), we obtain \( L^{-D}(\cdot) = L^{X_2-X_1}(\cdot) \equiv 0 \) as well as \( \hat{L}^D(\cdot) \equiv 0 \). Furthermore, by Corollary 2.6 of OUKNINE & RUTKOWSKI (1995), we obtain
\[
\hat{L}^{X_1\lor X_2}(t) = \int_0^t 1_{\{X_2(s) \leq 0\}} \, d\hat{L}^{X_1}(s) + \int_0^t 1_{\{X_1(s) < 0\}} \, d\hat{L}^{X_2}(s); \quad 0 \leq t < \infty.
\]
Combining these results with TANAKA’s formula, we obtain the dynamics of \( M(\cdot) := X_1(\cdot) \lor X_2(\cdot) \):
\[
dM(t) = 1_{\{X_1(t) \geq X_2(t)\}} \, dX_1(t) + 1_{\{X_1(t) < X_2(t)\}} \, dX_2(t) + dL^{X_1-X_2}(t)
\]
\[
= 1_{\{X_1(t) \geq X_2(t)\}} \left( \text{sgn}(X_1(t)) \, dU(t) + dV(t) + 2(2\alpha - 1) \, d\hat{L}^{X_1}(t) \right)
\]
\[
+ 1_{\{X_1(t) < X_2(t)\}} \left( \text{sgn}(X_2(t)) \, dU(t) + dV(t) + 2(2\alpha - 1) \, d\hat{L}^{X_2}(t) \right)
\]
\[
= \text{sgn}(M(t)) \, dU(t) + dV(t) + 2(2\alpha - 1) \, d\hat{L}^M(t); \quad 0 \leq t < \infty.
\]
In other words, each of the continuous semimartingales \( X_1(\cdot), \ X_2(\cdot) \) and \( M(\cdot) = X_1(\cdot) \lor X_2(\cdot) \) satisfies the equation (2.16); but uniqueness in the sense of the probability distribution holds for this equation, so all three processes have the same distribution. Since \( M(\cdot) \geq X_i(\cdot), \) this forces \( M(\cdot) = X_i(\cdot), \) \( i = 1, 2, \) thus pathwise uniqueness. By the theory of YAMADA and WATANABE (e.g., subsection 5.3.D in KARATZAS & SHREVE (1991)), the solution to (2.16) is therefore strong. The proof of Proposition 2.3 is complete.

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