High density terahertz frequency comb produced by coherent synchrotron radiation

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Frequency combs have enabled significant progress in frequency metrology and high-resolution spectroscopy extending the achievable resolution while increasing the signal-to-noise ratio. In its coherent mode, synchrotron radiation is accepted to provide an intense terahertz continuum covering a wide spectral range from about 0.1 to 1 THz. Using a dedicated heterodyne receiver, we reveal the purely discrete nature of this emission. A phase relationship between the light pulses leads to a powerful frequency comb spanning over one decade in frequency. The comb has a mode spacing of 846 kHz, a linewidth of about 200 Hz, a fractional precision of about $2 \times 10^{-10}$ and no frequency offset. The unprecedented potential of the comb for high-resolution spectroscopy is demonstrated by the accurate determination of pure rotation transitions of acetonitrile.
frequency combs (FC) have radically changed the landscape of optical instrumentation containing a vast number of narrow optical modes at a series of well-defined frequencies, which may be used as a highly accurate ruler over its spectral range. Initially developed in the visible and near-infrared spectral regions, the use of FC has been expanded to the mid-infrared, extreme ultra-violet, X-ray regions and significant effort is currently being dedicated to the generation of FC at THz frequencies. One solution based on converting a stabilized optical FC using a photoconductive terahertz emitter remains hampered by the low available THz power. Another approach is based on active mode locked THz quantum-cascade-lasers providing intense FC over a relatively limited spectral range.

As an alternative to laser based experiments and electronic generation, broadband THz radiation emitted from short electron bunches is now used for high-resolution spectroscopy and spectro-microscopy applications. In the so-called low-\(\alpha\) mode, intense THz power is obtained by reducing the electron bunch length to produce phase coherent synchrotron radiation (CSR). Following the first demonstration, such radiation has been developed in several synchrotron facilities over the world. It has provided unprecedented power in the region 0.1–1 THz, up to 4 orders of magnitude brighter than standard incoherent synchrotron radiation. Super-Radiance effects obtained from bunch-to-bunch interference further enhance specific frequencies (\(1/T_{\text{bunch}}\)) where \(T_{\text{bunch}}\) is the time interval between successive bunches of the THz CSR.

In this work, using a dedicated high-performance sub-THz heterodyne receiver (see Methods), we reveal that THz CSR is an intense, stable (in both frequency and amplitude), dense zero offset FC covering a wide spectral range 0.1–1 THz with no continuum background.

Results

Spectral characterization of CSR. In this work, the CSR spectral structure has been probed by a Bruker Fourier Transform Infrared (FTIR) instrument with an ultimate resolution of 30 MHz, as well as by the proposed heterodyne receiver shown in Fig. 1. The receiver, capable of operation in a large spectral region with a bandwidth of about 10 GHz, probes the synchrotron radiation with unprecedented spectral resolution, high signal-to-noise ratio and excellent frequency accuracy. Using both spectrometers, we are able to detect the fine spectral features of the CSR. In particular, we identified two main ultra-stable FC sequences in both regions studied (200 ± 10 GHz and 400 ± 10 GHz). The FC sequences are illustrated in Fig. 2 with a progressively finer frequency scale. The upper panel of this figure shows the CSR spectrum as recorded by the FTIR spectrometer at the highest possible spectral resolution (maximum optical path difference of 8.82 m, corresponding to a resolution of 30 MHz).

The full CSR range convoluted by FTIR response is shown in Fig. 2a. A detailed view ~600 GHz reveals a 352 MHz FC superimposed on the broad THz emission, Fig. 2b, as previously observed in other facilities. The lower panel shows the heterodyne analysis of the CSR. Figure 2c shows 3 GHz of heterodyne spectrum at 200 GHz, in which the 352 MHz FC originating from the bunch-to-bunch repetition is also clearly resolved. This FC is the signature of a degree of coherence among the electron bunches and is designated as super-radiance emission. Expanding the frequency scale further reveals a second FC composed of sharp teeth regularly spaced by 846 kHz, Fig. 2d. This comb is indeed related to the very stable revolution period of the electron bunches in the storage ring (1.18 s). It produces a spectrally dense THz FC with >10^8 components covering the THz range from 0.1 to 1 THz. The observed structure as shown in

**Figure 1 | CSR heterodyne detection schematic.** Short bunches of relativistic electrons circulating in the SOLEIL storage ring emit CSR in the 0.1–1 THz range. The bunch-to-bunch frequency \(1/T_{\text{bunch}}\) is 352 MHz, while the storage ring revolution frequency is 846 kHz (see Methods). The THz CSR pulses are mixed with the monochromatic radiation \(f_{\text{LO}}\) from a local oscillator (microwave synthesizer and multiplier) creating an electronic signal at the intermediate frequency (IF): \(f_{\text{IF}} = f_{\text{LO}} - f_{\text{CSR}}\) within the mixer bandwidth. The resulting signal is amplified by the IF amplifier A and analysed using a microwave spectrum analyser. The mixing produces a frequency downconversion of \(f_{\text{CSR}}\) into the microwave range, and an aliasing of the spectrum due to the superposition of the upper \(f_{\text{CSR}} > f_{\text{LO}}\) and the lower \(f_{\text{CSR}} < f_{\text{LO}}\) frequency bands.
Fig. 2d is composed of the lower and upper frequency band contributions as expected for heterodyne detection technique. Aliasing of the FC results in the observation of a doubled FC structure whose frequency separation is dependent on the local oscillator (LO) frequency. Shifting the LO frequency allowed the two components to be distinguished, here a separation of 100 kHz was selected, Fig. 2e. The detailed examination of any individual tooth, Fig. 2e inset, indicates an even narrower structure: typically, two components separated by 1.3 kHz, each with a full width at half maximum (FWHM) systematically smaller than 500 Hz. The substructure is interpreted as a consequence of the kHz range amplitude modulation caused by the low-frequency instabilities of the electron bunches.

**Offset determination.** In general, FC modes are described as multiples of a reference frequency \( f_R \) shifted by an offset \( f_0 \):

\[
f_{\text{FC}} = n \times f_R + f_0, \quad \text{where } n \text{ is an integer.}
\]

The offset frequency was assessed by measuring the revolution frequency with a counter (CNT 90, Pendulum). Over a period of 20 s, the revolution frequency is determined with a typical uncertainty of 30 mHz. The comb mode centre frequencies were determined by the spectral analysis of the intermediate frequency (IF) signal. In excess of 50 modes were examined, each being found to be an exact multiple of the revolution frequency with a confidence interval of 1 Hz. Unlike optical FC, the CSR FC is therefore considered to be offset free.

**Power measurement.** The heterodyne analysis allowed the power of an individual CSR line ~400 GHz to be estimated at −77 dBm (that is, 20 pW) (see Methods). The CSR spectral power density is 3 orders of magnitude greater than can be achieved from a classical dipolar antenna used in a context of THz time domain spectroscopy. The dynamic range of the amplitude of the FC teeth is systematically in excess of 35 dB with the strongest lines exhibiting over 70 dB.

**Pure rotational spectroscopy of acetonitrile.** We recorded the absorption spectrum of acetonitrile (CH\(_3\)CN) as the K series of this symmetric top molecule is particularly well suited to test the frequency accuracy and the dynamic range of the CSR FC. A 1.3-m-long absorption gas cell was inserted into the CSR path. Figure 3 shows a small section of the acetonitrile absorption spectrum obtained in 65 s using the CSR FC with the LO at 202 GHz. The FC gives rise to a discrete spectrum built from the harmonics of the repetition rate of 846 kHz (violet vertical lines).
The complete $K$ structure of the R(10) transition of CH$_3$CN is clearly resolved (violet profile, as explained in the associated content). This is in strong contrast to the simulated spectrum corresponding to the 30 MHz resolution of the IFS125 interferometer (red profile), only displaying a broad unresolved feature. In particular, the THz FC allows the doublet $K = 0$ and $K = 1$ separated by 3.9 MHz ($1.3 \times 10^{-4}$ cm$^{-1}$) to be resolved. The fit of the five individual absorption lines with Voigt profiles provides their central frequencies, which are in excellent agreement with reference data (r.m.s $< 80$ kHz) (see Methods). The accuracy obtained is comparable with frequency multiplication techniques currently employed in this frequency range.

**Discussion**

The measurements described in the previous section show clearly the discrete nature of the CSR. Moreover, the relatively high power together with the frequency stability of this source allows for room temperature heterodyne fast detection. Such results demonstrate that the phase relation between the emissions from bunches exists and is preserved over a long time to form a dense FC. In this sense, in the low alpha configuration, the operational mode of the synchrotron can be likened to ultrashort laser pulses where the first repetition rate is determined by the length of the ring resonator, and a second being defined by the bunch-to-bunch spacing. It is noteworthy that by using the same detection scheme, no discrete spectral components and no FC have been observed in standard synchrotron mode proving that for longer bunches of electrons there is no phase coherence between subsequent pulses, preventing the formation of a dense FC.

As shown by the quality of the THz acetonitrile absorption spectrum, the CSR FC is particularly suited to high-resolution gas phase THz spectroscopy as the mode spacing of 846 kHz is the same order of magnitude of the Doppler-limited linewidths of a typical rotational transition of a light polar compound at room temperature. This is not the case for many laser-generated THz combs with mode spacing typically $> 100$ MHz, requiring complex interleaving strategies to be adopted to obtain Doppler-limited rotational spectra.$^{6,15}$ Indeed the CSR FC composed of powerful, fine modes over an extended spectral range (0.1–1 THz) should outperform FTIR instruments in terms of resolution and analysis times. In comparison to electronic sources that offer an excellent resolution, the advantage of CSR FC is the spectral coverage of one decade, whereas a frequency multiplier may offer a tunability of $\sim 20\%$, numerous multiplier combinations are required to cover an octave of frequency. The excellent resolving power, frequency metrology and high dynamic range of the CSR FC are demonstrated by the high-resolution spectroscopy measurements undertaken using the heterodyne receiver initially constructed for the spectral analysis.
To harness the full potential of this FC, a solution must be sought to simultaneously access its entire frequency band (that is, over more than one decade of frequency). Mixing the CSR FC with an optical comb in a suitable photoconductor should yield the required information via a dual comb spectroscopy approach unlocking the advantages of the CSR\textsuperscript{11,13}.

The CSR FC characteristics revealed here clearly prompt the development of receivers for ultra-high resolution over broader frequency ranges. This is an enabling approach for new high-resolution THz time-resolved experiments. Among highly desirable new applications are ultrafast monitoring of reaction kinetics, ultra-high resolution atomic and molecular spectroscopy\textsuperscript{2,6}, and spatio-temporal dynamics of relativistic electron bunches.

### Methods

#### Coherent synchrotron radiation

The AILES beamline collecting the edge emission\textsuperscript{14} from the 2.75 GeV, third generation Synchrotron Radiation source SOLEIL, has an extended programme of research with synchrotron radiation in the THz spectral range. In particular, a special attention is paid into increasing the flux of CSR lines which should be around $\frac{E_{\text{tot}}}{C_{0}}$ for the required information via a dual comb spectroscopy approach seeking to simultaneously access its entire frequency band (that is, the THz spectral range. In particular, a special attention is paid into increasing the flux of CSR lines which should be around $\frac{E_{\text{tot}}}{C_{0}}$ for the required information via a dual comb spectroscopy approach.

\[ \frac{E_{\text{tot}}}{C_{0}} \quad \text{for the required information via a dual comb spectroscopy approach} \]

#### Heterodyne detection

To investigate the spectral structure of the THz CSR, we constructed a high sensitivity coherent electronic receiving\textsuperscript{18}. Heterodyne detection is very convenient for high-resolution spectral analysis with a high dynamic range. The operation is based on a frequency translation of the analysed spectra by electrical mixing between a known reference frequency (LO) and the signal to be detected. The mixing element is a Schottky-based sub-harmonic-mixer (SHM), fed by the LO at mW level (continuous wave), this power being required to operate the Schottky device in a non-linear regime. Two different SHM mixers WR5.1 (140–220 GHz) and WR2.2 (325–500 GHz) are used for operation at 200 GHz and 400 GHz, respectively. The LO is generated by an electronic frequency multiplier from a chain driven by a RF and Schoenholz synthesizer. This provides equivalent frequency multiplication factors of \(12\times\) and \(12\times\), respectively. The IF signal is amplified with two broad-band amplifiers, providing a total gain of 45 dB over a bandwidth of 500 MHz. The downconverted CSR signal is analysed by a 3.6 GHz Agilent MXA Signal Analyzer N9020A.

#### Power measurements

From IF signal spectral analysis, the individual CSR lines located around 400 GHz display a typical signal strength of \(-40\) dBm. Taking into account the 45 dB IF gain and a 8 dB conversion loss of SHM, the corresponding power of CSR lines should be around \(-77\) dBm (20 pW). In integrated THz power of the CSR was measured to be 60.5 mW by a Gentec THz SB-DLZ pyroelectric detector working at room temperature. The beam was filtered with a black polyethylene film to avoid any effect of visible or mid-infrared radiation. The ratio of the integrated power to the power of one line is then equal to about \(3 \times 10^6\). This gives a rough estimation of the number of lines assuming that all lines of the FC have the same intensity. This value is compatible with the fact that in a THz range there should be about \(12 \times 10^6\) lines separated by 846 kHz. The spectral power density is hence calculated to be 23.6 mW GHz\(^{-1}\).

#### Spectroscopic analysis on acetonitrile

The IF spectra are processed to isolate the lower and upper frequency band contributions. Once all the peaks in the spectra have been identified, the FC spacing of 846 kHz is used to distinguish each FC contribution. In Fig. 2, the normalized absorption of CH\(_3\)CN measured with the THz FC was simulated with Voigt profiles. The Doppler linewidth (Gaussian contributions) of the five components has been fixed to a calculated value at room temperature of 0.4 MHz (FWHM). The centre frequencies and the collisional linewidths (Lorentzian contributions) were fitted for \(K=0-5\) components of the R(10) transitions. In Table 1, the fitted frequencies \(f_{\text{fit}}\) are compared with the frequencies \(f_{\text{calc}}\) listed in the JPL database\textsuperscript{19} and calculated from the latest combined fit reported by Müller et al.\textsuperscript{20}. The r.m.s. deviation of the fitted frequencies \(f_{\text{fit}}\) from the calculated \(f_{\text{calc}}\) values is evaluated to be 79 kHz. The collisional linewidths were fitted to 3.4 MHz (FWHM) for the five components (no \(K\) dependence of the R(10) self-broadening was observed). At a pressure of 50 mbar a collisional linewidth of \(\sim 4.5\) MHz (FWHM) may be estimated for the R(10) transitions from CH\(_3\)CN self-broadening measurements\textsuperscript{21,22}.

### Table 1 fitted frequencies compared with the frequencies listed in the JPL database\textsuperscript{19}.

| \(K\) | \(f_{\text{fit}}\) (MHz) | \(f_{\text{calc}}\) (MHz) | \(100 \times (f_{\text{fit}} - f_{\text{calc}})/f_{\text{calc}}\) |
|------|------------------|------------------|------------------|
| 0    | 202,355.39       | 202,355.51       | 0.12             |
| 1    | 202,351.73       | 202,351.61       | 0.12             |
| 2    | 202,339.88       | 202,339.92       | 0.04             |
| 3    | 202,320.42       | 202,320.44       | 0.02             |
| 4    | 202,293.19       | 202,293.18       | 0.01             |

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O.P. and G.M. wrote the synchrotron beamtime proposal and supervised the project. Experiments were performed by S.T., O.P., J.-F.L., G.D., F.H. and G.M. All the authors contributed equally to the analysis of the results and to the preparation of the manuscript.

Additional information
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