CUGLIANDOLO-KURCHAN EQUATIONS FOR DYNAMICS OF SPIN-GLASSES.

GERARD BEN AROUS, AMIR DEMBO, AND ALICE GUIONNET

Abstract. We study the Langevin dynamics for the family of spherical $p$-spin disordered mean-field models of statistical physics. We prove that in the limit of system size $N$ approaching infinity, the empirical state correlation and integrated response functions for these $N$-dimensional coupled diffusions converge almost surely and uniformly in time, to the non-random unique strong solution of a pair of explicit non-linear integro-differential equations, first introduced by Cugliandolo and Kurchan.

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1. Introduction and main results

Markovian dynamics with random interactions can produce very complex phase transitions, and fascinating long time behaviors, for strong disorder (or low temperature). The physics literature has shown that dynamics of mean-field spin glasses is a very good field to get an accurate sample of possible and generic long time phenomena (as aging, memory, rejuvenation, failure of the Fluctuation-Dissipation theorem, see [8] for a good survey).

This class of problems can be roughly described as follows. Let $\Gamma$ (a compact metric space) be the state space for spins and $\nu$ be a probability measure on $\Gamma$. Typically, in the discrete or Ising spins context $\Gamma = \{-1, 1\}$ and $\nu = 1/2(\delta_1 + \delta_{-1})$. In the continuous or soft spin context, $\Gamma = I$ a compact interval of the real line and $\nu(dx) = Z^{-1}e^{-U(x)}dx$, where $U(x)$ is the "one-body potential".

For each configuration of the spin system, i.e. for each $x = (x_1, ..., x_N) \in \Gamma^N$ one defines a random Hamiltonian, $H^N_J(x)$, as a function of the configuration $x$ and of an exterior source of randomness $J$, i.e. a random variable defined on another probability space. The Gibbs measure at inverse temperature $\beta$ is then defined on the configuration space $\Gamma^N$ by

$$\mu^N_{\beta, J}(dx) = \exp(-\beta H^N_J(x))\nu(dx)/Z^N_J$$

The statics problem amounts to understanding the large $N$ behavior of these measures for various classes of random Hamiltonians ([21] is a recent and beautiful book on the mathematical results pertaining to these equilibrium problems). The dynamics question consists of understanding the behavior of Markovian processes on the configuration space $\Gamma^N$, for which the Gibbs measure is invariant and even reversible, in the limit of large systems (large $N$) and long times, either when the
randomness $\mathbf{J}$ is fixed (the quenched case) or when it is averaged (often called the annealed case, in
the mathematics literature, but not in the physics papers).

These dynamics are typically Glauber dynamics for the discrete spin setting, or Langevin dynamics
for continuous spins. Defining precisely what we mean here by large system size and long time is a
very important question, and very different results can be expected (and sometimes proved, see [3]
and references therein), for various time-scales as functions of the size of the system. We will restrict
ourselves to the case where the system size is first taken to infinity, i.e for the shortest possible long
time scales, much too short typically to allow any escape from meta-stable states (as opposed to the
situation in [4] for instance). The first step is then to derive limiting equations for various quantities,
when $N$ tends to $\infty$. The second step is to understand the large time behavior of these limiting
macroscopic equations. Let us be more specific by describing one of the main initial objectives of
the theory, which is the long time behavior of the Langevin dynamics of the Sherrington-Kirkpatrick
model and its generalization, the $p$-spin model. In the SK model, either for discrete or for continuous
spins, the Hamiltonian is given by:

$$H^N_3(\mathbf{x}) = \sum_{1 \leq i,j \leq N} J_{ij} x^i x^j,$$

where the randomness is due to the couplings $J_{ij}$ which are assumed to be i.i.d Gaussian centered,
of variance $N^{-1}$ in case $i \neq j$ and $2N^{-1}$ in case $i = j$. In the $p$-spins model the Hamiltonian contains
interaction between subsets of spins of size $p$, that is,

$$(1.1) \quad H^N_3(\mathbf{x}) = \sum_{1 \leq i_1 < ... < i_p \leq N} J_{i_1...i_p} x^{i_1} ... x^{i_p},$$

where the couplings $J_{i_1...i_p}$ are assumed to be i.i.d. Gaussian centered and of variance of $O(N^{-(p-1)})$.
The initial SK model corresponds of course to the case $p = 2$.

Propagation of chaos for dynamics of the SK model, or equivalently the large $N$ limit for the
behavior of the empirical measure $\frac{1}{N} \sum_{i=1}^{N} \delta_{x_i(t)}$ has been obtained some time ago both for continuous
spins (see [19] [20], in the physics literature and [6] [7] [15] for a mathematical treatment) or discrete
ones (see [14]). The limiting equations (the so called self-consistent single-spin dynamics) are very
complex and have resisted so far all attempts to understand their long time behavior. This is due
in part to the fact that the empirical measure is a much too rich object. Finding an autonomous
system of tractable equations for a proper well chosen set of lower dimensional quantities is still a
very open question, even in the physics literature.

But a large range of interesting and related models have been recently analyzed more successfully
in the physics literature, i.e the spherical $p$-spin models (see [8] [10]). The spherical version of the
spin glass models consists of a classical simplification, that is, replacing the product structure of the
configuration space $\Gamma_N$ by the sphere $S^{N-1}(\sqrt{N})$ of radius $\sqrt{N}$ in $\mathbb{R}^N$, in effect, imposing on the
configuration $\mathbf{x}$ a hard constraint $\frac{1}{N} \sum_{i=1}^{N} x_i^2 = 1$. The spherical $p$-spin Gibbs measure is thus the
probability measure on the sphere $S^{N-1}(\sqrt{N})$ given by

$$\mu^N_{\beta,3}(d\mathbf{x}) = \exp(-\beta H^N_3(\mathbf{x})) \nu_N(d\mathbf{x})/Z^N_3,$$

where the Hamiltonian is given by (1.1) and the measure $\nu_N$ is the uniform measure on the sphere
$S^{N-1}(\sqrt{N})$. One can also study a very similar problem by replacing the hard spherical constraint by
a soft one, i.e by replacing the uniform measure $\nu_N$ on the sphere $S^{N-1}(\sqrt{N})$ by a measure on $\mathbb{R}^N$

$$\nu_N(d\mathbf{x}) = \exp\left(-N f\left(\frac{1}{N} \sum_{i=1}^{N} x_i^2\right)\right) dx/Z^N,$$
where $f$ is a smooth function growing fast enough at infinity. This study has been done successfully (see \cite{5}, and \cite{11}) in the case where $p = 2$, the spherical SK model. There one could obtain a very complete description of the limiting dynamics using only one quantity, the empirical state correlation function

$$C_N(s, t) = \frac{1}{N} \sum_{i=1}^{N} x^i(s)x^i(t).$$

Indeed, this quantity was shown to have a non-random limit $C(s, t)$, when $N$ tends to $\infty$, which satisfies an autonomous integro-differential equation. This is a rather simple setting, when $p = 2$, since this quadratic case can use the well known tools pertaining to the study of spectra of GOE random matrices, after diagonalization of the random matrix $(J_{ij})$. The results, though showing an aging phenomenon, are very far from what is expected for the true (i.e non spherical) SK model.

The case of $p > 2$ is a completely different story. The physics literature, mainly Cugliandolo and Kurchan (see \cite{8,10,12}), has given a very rich picture of the behavior of the Langevin dynamics for these spherical $p$-spin models, which are believed to mimic in certain cases the behavior of the dynamics of the true SK model. The main innovation is the fact that the description of the limiting dynamics relies now on two objects, the empirical correlation function and the response function (and not only one as in the spherical $p = 2$ case). Our work establishes rigorously the asymptotic validity of these Cugliandolo-Kurchan equations (see \cite{12} for the physical derivation of these). We shall devote a future work to the fascinating question of understanding the large time behavior of the solution to these equations (see \cite{8} or \cite{10} to see what is expected).

Some of the needed mathematical tools for attacking this problem can be found in \cite{16}, as well as an intriguing link with equations arising naturally in non-commutative probability.

We now describe more precisely our result and then the structure of this paper.

Fixing a positive integer $N$ (denoting the system size), we consider the mean field random Hamiltonian

$$H^N_j(x) = -2 \sum_{p=1}^{m} \frac{a_p}{p!} \sum_{1 \leq i_1, \ldots, i_p \leq N} J_{i_1 \ldots i_p} x^{i_1} \ldots x^{i_p},$$

where $m \geq 2$, the state variable is $x = (x^1, \ldots, x^N) \in \mathbb{R}^N$, and the disorder parameters $J_{i_1 \ldots i_p} = J_{(i_1, \ldots, i_p)}$ are independent (modulo the permutation of the indices) centered Gaussian variables. The variance of $J_{i_1 \ldots i_p}$ is $c(|\{i_1, \ldots, i_p\}|)N^{-p+1}$, where

$$c(|\{i_1, \ldots, i_p\}|) = \prod_k l_k!,$$

and $(l_1, l_2, \ldots)$ are the multiplicities of the different elements of the set $\{i_1, \ldots, i_p\}$ (for example, $c = 1$ when $i_j \neq i'_j$ for any $j \neq j'$, while $c = p!$ when all $i_j$ values are the same).

Let $f$ be a differentiable function on $\mathbb{R}_+$ with $f'$ locally Lipschitz, such that

$$\sup_{\rho \geq 0} |f'(\rho)|(1 + \rho)^{-r} < \infty$$

for some $r < \infty$, and for some $A, \delta > 0$,

$$\inf_{\rho \geq 0} \{f'(\rho) - A\rho^{m/2 + \delta - 1}\} > -\infty$$

(typically, $f(\rho) = \kappa(\rho - 1)^r$ for some $r > m/2$ and $\kappa \gg 1$). At temperature $\beta^{-1} > 0$, the soft spherical version of the system corresponds to the equilibrium probability measure $\mu^N_{\beta,J}$ on $\mathbb{R}^N$ whose density
norm for the Brownian motion $B$ equipped with the natural Euclidean norms for the finite dimensional parts, i.e.
\( f(x) = \sum \frac{a_p}{(p - 1)!} \sum_{1 \leq i_1, \ldots, i_p \leq N} J_{i_1 \ldots i_p} x^{i_1} \ldots x^{i_p}, \)

for $i = 1, \ldots, N$.

In Proposition 2.1 we prove that for a.e. disorder $J$, initial condition $x_0$ and Brownian path $B$, there exists a unique strong solution of (1.6) for all $t \geq 0$, whose law we denote by $\mathbb{P}^N_{\beta, x_0, J}$.

We are interested in the time evolution of the empirical state correlation function
\[
C_N(s, t) := \frac{1}{N} \sum_{i=1}^{N} x_i^s x_i^t,
\]
and that of the empirical integrated response function
\[
\chi_N(s, t) := \frac{1}{N} \sum_{i=1}^{N} x_i^s B_i^t,
\]
under the quenched law $\mathbb{P}^N_{\beta, x_0, J}$, as the system size $N \to \infty$. Note that since $\{\chi_N(s, t), 0 \leq t \leq s \leq T\}$ is not determined by $N^{-1} \sum \delta_{x_i^0, x_j^0}$, the strategy by which the limiting equations in [6, 7] are derived, does not apply directly in our case.

We assume hereafter that the initial condition $x_0$ is independent of the disorder $J$, and the limit
\[
\lim_{N \to \infty} \mathbb{E} C_N(0, 0) = C(0, 0),
\]
exists, and is a finite. Further, we assume that the tail probabilities $\mathbb{P}(|C_N(0, 0) - C(0, 0)| > x)$ decay exponentially fast in $N$ (so the convergence $C_N(0, 0) \to C(0, 0)$ holds almost surely), and that for each $k < \infty$, the sequence $N \to \mathbb{E}[C_N(0, 0)^k]$ is uniformly bounded.

To be more specific, we consider hereafter the product probability space $\mathcal{E}_N = \mathbb{R}^N \times \mathbb{R}^{d(N, m)} \times C([0, T], \mathbb{R}^N)$ (here $T$ is a fixed time and $d(N, m)$ is the dimension of the space of the interactions $J$), equipped with the natural Euclidean norms for the finite dimensional parts, i.e. $(x_0, J)$, and the sup-norm for the Brownian motion $B$. The space $\mathcal{E}_N$ is endowed with the product probability measure $\mathbb{P} = \mu_N \otimes \gamma_N \otimes P_N$, where $\mu_N$ denotes the distribution of $x_0$, $\gamma_N$ is the (Gaussian) distribution of the coupling constants $J$, and $P_N$ is the distribution of the $N$-dimensional Brownian motion.

**Hypothesis 1.1.** For $(x_0, J, B) \in \mathcal{E}_N$ we introduce the norms
\[
\|(x_0, J, B)\|_2^2 = \sum_{i=1}^{N} (x_0^i)^2 + \sum_{p=1}^{m} \sum_{1 \leq i_1 \ldots i_p \leq N} \left( N^{p-1} J_{i_1 \ldots i_p} \right)^2 + \sup_{0 \leq t \leq T} \sum_{i=1}^{N} (B_t^i)^2.
\]
We shall assume that \( \mu_N \) is such that the following concentration of measure property holds on \( \mathcal{E}_N \): there exists two finite positive constants \( C \) and \( \alpha \), independent on \( N \), such that, if \( V \) is a Lipschitz function on \( \mathcal{E}_N \), with Lipschitz constant \( K \), then for all \( \rho > 0 \),

\[
\mu_N \otimes \gamma_N \otimes P_N |V - \mathbb{E}[V]| \geq \rho \leq C^{-1} \exp \left( - \frac{\rho^\alpha}{K} \right).
\]

The above concentration inequality holds for any Lipschitz function \( V \) that does not depend on \( x_0 \), since the concentration of measure property holds for the Gaussian measures \( \gamma_N \otimes P_N \), with \( \alpha = 2 \) (c.f. \[2\]). Unfortunately, assuming that the concentration of measure property holds for the measures \( \mu_N \) on \( \mathbb{R}^N \) does not assure the concentration of measure for the product measure \( \mu_N \otimes \gamma_N \otimes P_N \). Hence, we shall have to assume a property which implies the concentration inequality of Hypothesis 1.1 and further tensorizes (see \[2\] for a more thorough discussion). For example, if \( \mu_N \) satisfy the Poincaré inequality uniformly in \( N \), then the product measures \( \mu_N \otimes \gamma_N \otimes P_N \) also satisfy the Poincaré inequality uniformly in \( N \), since the Gaussian measure \( \gamma_N \otimes P_N \) does. The required uniform in \( N \) concentration property is then satisfied, with \( \alpha = 1 \) (c.f. \[12\]).

Our main result is the proof that as \( N \to \infty \) the functions \( C_N(s,t) \) and \( \chi_N(s,t) \) converge to non-random functions \( C(s,t) \) and \( \chi(s,t) \), that are characterized by the following theorem.

**Theorem 1.2.** Let \( \psi(r) = \nu'(r) + r \nu''(r) \) and

\[
(1.11) \quad \nu(r) := \sum_{p=1}^{m} \frac{a_p^2}{p!} r^p.
\]

Suppose \( \mu_N \) satisfies hypothesis 1.1. Fixing any \( T < \infty \), as \( N \to \infty \) the random functions \( C_N \) and \( \chi_N \) converge uniformly on \([0,T]^2\), almost surely and in \( L_p \) with respect to \( x_0 \), \( J \) and \( B \) to non-random functions \( \chi(s,t) = \int_0^t R(s,u)du \) and \( C(s,t) = C(t,s) \). Further, \( R(s,t) = 0 \) for \( t > s \), \( R(s,s) = 1 \), and for \( s > t \) the absolutely continuous functions \( C, R \) and \( K(s) = C(s,s) \) are the unique solution in the space of bounded, continuous functions, of the integro-differential equations

\[
\begin{align*}
(1.32) R(s,t) &= -f'(K(s))R(s,t) + \beta^2 \int_0^t R(u,t)R(u,s)\nu''(C(u))du, \\
(1.33) C(s,t) &= -f'(K(s))C(s,t) + \beta^2 \int_0^t C(u,t)R(u,s)\nu''(C(u))du + \beta^2 \int_0^t \nu'(C(u))R(t,u)du, \\
(1.34) K(s) &= -2f'(K(s))K(s) + 1 + \beta^2 \int_0^t \psi(C(u))R(s,u)du,
\end{align*}
\]

where the initial condition \( K(0) = C(0,0) \) is determined by 1.10.

We note in passing that for \( p = 2 \), i.e. \( \nu(r) = r^2/2 \), we get from 1.12 the autonomous equation

\[
\partial_s H(s,t) = \beta^2 \int_0^s H(u,t)H(s,u)du,
\]

for \( H(s,t) = R(s,t) \exp(\int_t^s f'(K(u))du) \), whose unique solution is the Laplace transform of the semicircle probability measure, evaluated at \( \beta(s-t) \). Plugging this expression in 1.13 and 1.14, upon scaling both the function \( f \) and time by a factor \( \beta \), we recover the limiting equation of \([5\] (2.16)] after some integrations by parts. Further, setting \( K(s) = 1 \) and \( \partial_s K(s) = 0 \) in 1.14, while replacing \( f'(K(s)) \) in 1.12–1.14 by a time varying constant \( z(s) \), corresponds to the hard spherical constraint of 12. Indeed, the limiting equations of 12 are thus recovered.

Note that the constant \( \beta \) can be embedded into \( \{a_p\} \) resulting with \( \beta G^2(\cdot) \to G^2(\cdot) \) and then having \( \beta = 1 \) in the stochastic differential system 1.6. Adopting this convention, we thus take hereafter \( \beta = 1 \). It is trivial to check that the explicit dependence of 1.12–1.14 on \( \beta \) is indeed as stated, i.e., with each appearance of \( \nu'() \), \( \nu''() \) (and \( \psi() \)) multiplied by \( \beta^2 \).
The empirical quantities $K_N(s) := C_N(s, s)$ and
\begin{equation}
A_N(s, t) := \frac{1}{N} \sum_{i=1}^{N} G^i(x_s) x_t^i, \quad F_N(s, t) := \frac{1}{N} \sum_{i=1}^{N} G^i(x_s) B_t^i,
\end{equation}
play a key role in the derivation of Theorem 1.2. Indeed, with
\begin{equation}
D_N(s, t) := -f'(\mathbb{E}(K_N(t))) C_N(s, t) + A_N(t, s), \quad E_N(s, t) := -f'(\mathbb{E}(K_N(s))) \chi_N(s, t) + F_N(s, t),
\end{equation}
the key step of the proof of Theorem 1.2 is summarized by

**Proposition 1.3.** Fixing any $T < \infty$, in case $\beta = 1$, any limit point of the sequence $(\mathbb{E}C_N, \mathbb{E}\chi_N, \mathbb{E}D_N, \mathbb{E}E_N)$ with respect to uniform convergence on $[0, T]^2$, satisfies the integral equations
\begin{align}
(1.17) \quad & C(s, 0) + \chi(s, t) + \int_0^t D(s, u) du, \\
(1.18) \quad & s \wedge t + \int_0^s E(u, t) du, \\
D(s, t) &= -f'(C(t, t)) C(t, s) - \int_0^{t \wedge s} \nu'(C(t, u)) D(s, u) du - \int_0^{t \wedge s} C(s, u) \nu''(C(t, u)) D(t, u) du \\
E(s, t) &= -f'(C(s, s)) \chi(s, t) - \int_0^s \nu'(C(s, u)) E(u, t) du - \int_0^s \chi(u, t) \nu''(C(s, u)) D(s, u) du \\
(1.19) \quad & + \chi(s, t) \nu'(C(s, s)) - \int_0^{t \wedge s} \nu'(C(s, u)) du,
\end{align}
in the space of bounded continuous functions on $[0, T]^2$ subject to the symmetry condition $C(s, t) = C(t, s)$ and the boundary conditions $E(s, 0) = 0$ for all $s$, and $E(s, t) = E(s, s)$ for all $t \geq s$.

We next detail the organization of the paper, and hence, that of the proof of Theorem 1.2.

In Section 2, we prove the existence of strong solutions $x_t$ for the Stochastic Differential System (in short SDS) given in (1.16), for any $N < \infty$ (see Proposition 2.1). We then prove that the functions $A_N, F_N, \chi_N$ and $C_N$ associated with these solutions have uniformly bounded (in $N$) finite moments of all order, and form pre-compact sequences with respect to uniform convergence on compacts both almost surely and in the mean (see Proposition 2.3). Applying the “localized concentration of measure” of Lemma 3.2, we complete the preliminary analysis of the finite size SDS by proving in Proposition 2.4 that as $N \to \infty$ each of the four functions $A_N, F_N, \chi_N$ and $C_N$ “self-averages”, namely, concentrates around its mean. These results rely on the bounding in Appendix B norms of the disorder $J$ that are expressed as the supremum of certain Gaussian fields.

The proof of Proposition 1.3, namely, that each limit point of $(\mathbb{E}C_N, \mathbb{E}\chi_N, \mathbb{E}D_N, \mathbb{E}E_N)$ must satisfy (1.17)–(1.20), is the subject of Section 3. Of these equations, upon multiplying the integrated form (3.1) of our SDS by $x_t^i$ or by $B_t^i$, then averaging over $i$ and the probability space, both (1.17) and (1.18) are immediate consequences of self-averaging. The crux of the proof is thus Proposition 3.1, where we show that as $N \to \infty$, both $\mathbb{E}A_N$ and $\mathbb{E}F_N$ are well approximated by certain combination of our four functions, thereby leading to (1.19) and (1.20). The emergence of $\nu'(C(s, u))$ and $\nu''(C(s, u))$ in the latter pair of equations, and hence in Theorem 1.2, is a consequence of the structure of the covariance kernel $k_N(x) = \mathbb{E}[G^i(x_t)G^j(x_s)]$, obtained by integrating over the disorder parameters $J$ assuming their independence of the frozen path $x_t$ (see Lemma 3.2). However, the main difficulty of the proof is the intricate dependence between $J$ and $\{x_t, 0 \leq t \leq T\}$. Taking full advantage of the Gaussian law of $J$ and the Brownian law of $B$, this difficulty is dealt with by combining the
Itô’s calculus identities of Appendix A with Girsanov’s theorem and the resulting Gaussian change of measure identities that are derived in Appendix B. This approach succeeds in deriving a “closed” system of finitely many limiting equations thanks to the fact that apart from our self-averaging global quantities, the kernel $k_t$ is a quadratic form of $x$. Note that certain Hamiltonians other than (1.2) also have such a property, hence are amenable to a similar treatment (one such example is $H^N_1(x/\|x\|)$ for $H^N_1(\cdot)$ of (1.2)).

Section 4 mostly deals with analytic considerations. Its starting point is Lemma 4.1 showing that any solution of (1.17)–(1.20) is sufficiently differentiable to give rise to a solution of (1.12)–(1.14). This is followed by Proposition 4.2, establishing the uniqueness of the latter system of equations any solution of (1.12)–(1.14), it is a probability measure on $C(\mathbb{R}^+, \mathbb{R}^N)$ which we denote $\mathbb{P}^N_{x_0,J}$. Further, with

\begin{equation}
(2.1) \quad \|J\|^N = \max_{1 \leq p \leq m} \sup_{1 \leq i \leq 1, 1 \leq i \leq p} \left[ \sqrt{N}^{-1} \sum_{1 \leq i \leq N, 1 \leq k \leq p} N^{\frac{i-1}{2}} J_{i_1 \cdots i_p} u_{i_1}^1 \cdots u_{i_p}^p \right]
\end{equation}

we have for $\delta > 0$ of (1.2), $q := m/(2\delta) + 1$, some $\kappa < \infty$, all $N$, $z > 0$, $J$, and $x_0$, that

\begin{equation}
(2.2) \quad \mathbb{P}^N_{x_0,J} \left( \sup_{t \in \mathbb{R}^+} K_N(t) \geq K_N(0) + \kappa \left( 1 + \|J\|^N_\infty \right)^q + z \right) \leq e^{-zN}.
\end{equation}

Consequently, for any $L > 0$, there exists $z = z(L) < \infty$ such that

\begin{equation}
(2.3) \quad \mathbb{P} \left( \sup_{t \in \mathbb{R}^+} K_N(t) \geq z \right) \leq e^{-LN}.
\end{equation}

**Proposition 2.1.** Assume that $f'$ is locally Lipschitz, satisfying (1.3). Then, for any $N \in \mathbb{N}$, almost any $J$, initial condition $x_0$ and Brownian path $B$, there exists a unique strong solution to (1.6). This solution is also unique in law for almost any $J$, and $x_0$, it is a probability measure on $C(\mathbb{R}^+, \mathbb{R}^N)$ which we denote $\mathbb{P}^N_{x_0,J}$. Further, with

\begin{equation}
(2.4) \quad \sum_{M=1}^{\infty} \mathbb{P} (\tau_M \leq T) < \infty.
\end{equation}

\textbf{Proof:} For every $M > 0$ we introduce a bounded globally Lipschitz function $\phi_M$ on $\mathbb{R}^N$ which we choose such that $\phi_M(x) = x$ when $\|x\| \leq \sqrt{NM}$, and then consider the truncated drift $b^M(u) = (b^M_1(u), \ldots, b^M_N(u))$ given by $b^M_i(u) = G^i(\phi_M(u)) - f'(N^{-1}\|u\|^2 \wedge M\hat{u})$.

Since $f'$ is locally Lipschitz, and since $\|J\|^N_\infty$ is finite almost surely for all $p$ and $N$, it is thus clear that the drift $b^M(u)$ is globally Lipschitz. The existence and uniqueness of a square-integrable strong solution $u^{(M)}$ for the SDS

\begin{equation}
du^M_t = b^M_t(u_t)dt + dB^M_t \quad \text{is thus standard (for example, see [17 Theorems 5.2.5, 5.2.9])}. \quad \text{With } u^{(M)} \text{ defined for all } M \text{ on the same probability space and filtration, consider the stopping times } \tau_M = \inf \{ t : \|u^{(M)}_t\| \geq \sqrt{NM} \}.
\end{equation}

Note that $u^{(M)}$ is the unique strong solution of (1.6) for $t \in [0, \tau_M]$, with $\tau_M$ a non-decreasing sequence. By the Borel-Cantelli lemma, it suffices for the existence of a unique strong solution $u = \lim_{M \to \infty} u^{(M)}$ of the SDS (1.6) in $[0, T]$, to show that

\begin{equation}
(2.4) \quad \sum_{M=1}^{\infty} \mathbb{P} (\tau_M \leq T) < \infty.
\end{equation}
To this end, fix $M$ and let $x_t = u^{(M)}_{s \wedge t}$ and $Z_s = 2N^{-1} \sum_{i=1}^N f^{s \wedge t} x_i dB_i$. Applying Itô’s formula for $C_N(t) = N^{-1}|x_t|^2$ we see that

$$C_N(s) \leq C_N(0) + 2 \sum_{p=1}^{m} \frac{a_p|J|^N}{(p-1)!} \int_0^{s \wedge t} C_N(t) \frac{r}{r^q} dt - 2 \int_0^{s \wedge t} f'(C_N(t))C_N(t) dt + Z_s + s \wedge \tau_M.$$  

(2.5)

Since $x^\frac{1}{2}f'(x) \to \infty$, it follows from (2.5) that there is an almost surely finite constant $c(||J||_\infty^N)$, independent of $M$, such that

$$C_N(s) \leq C_N(0) + c(||J||_\infty^N)s + Z_s$$  

(2.6)

As the quadratic variation of the martingale $Z_s$ is $(4/N) \int_0^{s \wedge t} C_N(t) dt \leq 4sN^{-1}M$, applying Doob’s inequality (c.f. [17, Theorem 3.8, p. 13]) for the exponential martingale $e^{M/s} = \exp(\lambda Z_s - 2(\lambda^2/N) \int_0^{s \wedge t} C_N(t) dt)$ (with respect to the filtration $\mathcal{H}_t$ of $\mathcal{B}_t$), yields that

$$\mathbb{P} \left( \sup_{s \leq t} \{Z_s - 2 \int_0^s C_N(t) dt\} \geq z \right) \leq \mathbb{P} \left( \sup_{s \leq t} L_s^N \geq e^{zN} \right) \leq e^{-zN},$$  

(2.7)

for any $z > 0$. Therefore, (2.6) shows that with probability greater than $1 - e^{-zN}$,

$$C_N(s \wedge \tau_M) \leq C_N(0) + c(||J||_\infty^N)T + z + 2 \int_0^{s \wedge t} C_N(t) dt,$$

for all $s \leq T$, and by Gronwall’s lemma then also

$$\sup_{t \leq T} N^{-1}|u^{(M)}_{s \wedge t}|^2 \leq |C_N(0) + c(||J||_\infty^N)T + z|e^{2T}.$$  

(2.8)

Setting $z = M/3$, for large enough $M$ (depending of $N, J, x_0$ and $T$ which are fixed here), the right-side of (2.8) is at most $M/2$, resulting with

$$\mathbb{P}(\tau_M \leq T) \leq e^{-MN/3},$$

and hence with (2.4). We also have weak uniqueness of our solutions for almost all $J$ since the restriction of any weak solution to the stopped $\sigma$-field $\mathcal{H}_{\tau_M}$ for the filtration $\mathcal{H}_t$ of $\mathcal{B}_t$ is unique. We denote this unique weak solution of (1.6) by $N_{x_0, J}$.

Turning to the proof of (2.2), by (1.5), for any $c > 0$ there exists $\kappa < \infty$ such that for all $r, x \geq 0$,

$$2[f'(x)x - r \sum_{p=1}^{m} \frac{a_p x^p}{(p-1)!}] - 1 \geq cx - \kappa(1 + r)^q.$$

Taking $r = ||J||_\infty^N$, we see that by (2.5), for all $N$ and $s \geq 0$,

$$C_N(s \wedge \tau_M) \leq C_N(0) - \int_0^{s \wedge t} [cC_N(t) - \kappa(1 + ||J||_\infty^N)^q] dt + Z_s,$$

where $(Z_s)_{s \geq 0}$ is a martingale with bracket $(4N^{-1} \int_0^{s \wedge t} C_N(t) dt, s \geq 0)$.

By Doob’s inequality (2.7), with probability at least $1 - e^{-zN}$,

$$\sup_{u \leq s \wedge \tau_M} Z_u \leq 2 \int_0^{s \wedge \tau_M} C_N(t) dt + z,$$

for all $s \geq 0$. Setting $c = 3$ we then have that

$$C_N(s \wedge \tau_M) \leq C_N(0) + z - \int_0^{s \wedge \tau_M} C_N(t) dt + \kappa(1 + ||J||_\infty^N)^q(s \wedge \tau_M),$$

(2.9)
so that by Gronwall’s lemma,
\[
C_N(s \wedge \tau_M) \leq e^{-s \wedge \tau_M} (C_N(0) + z) + \kappa (1 + \|J\|_\infty)^q \int_0^{s \wedge \tau_M} e^{-t} \, dt
\]
from which the conclusion (2.2) is obtained by considering Gronwall’s lemma, (B.7) on the corresponding probabilities for \( |J|_N \) we thus get also the bounds of (2.3).

Our next lemma provides bounds on \( G^i(x) \) of (1.7) which we often use en-route to the uniform boundedness of moments, pre-compactness, and concentration around the mean of the functions \( A_N, F_N, \chi_N, \) and \( C_N \).

**Lemma 2.2.** There exists a constant \( c < \infty \) such that, for all \( N < \infty \) and every \( x, \bar{x} \in \mathbb{R}^N \),
\[
\left( \sum_{i=1}^N (G^i(x) - G^i(\bar{x}))^2 \right)^{\frac{1}{2}} \leq c ||J||_\infty [1 + \left( \frac{1}{N} ||x||^2 \right)^{\frac{m-2}{2}}] [1 + \left( \frac{1}{N} ||\bar{x}||^2 \right)^{\frac{m-2}{2}}] ||x - \bar{x}||,
\]
and in particular,
\[
\left( \frac{1}{N} \sum_{i=1}^N G^i(x)^2 \right)^{\frac{1}{2}} \leq c ||J||_\infty [1 + \left( \frac{1}{N} ||x||^2 \right)^{\frac{m-2}{2}}].
\]

**Proof:** Fixing \( N, x, \bar{x} \) and \( y \in \mathbb{R}^N \), note the telescoping sum
\[
\sum_{i=1}^N (G^i(x) - G^i(\bar{x})) y^i = \sum_{p=1}^m \frac{a_p}{(p-1)!} \sum_{l=1}^{p-1} \sum_{i_k} J_{i_1 \cdots i_p} \left( \prod_{l=1}^{l-1} x^{i_l} \right) \left( \prod_{h=l+1}^{p-1} \bar{x}^{i_h} \right) \prod_{l=1}^p y^{i_l}
\]
Recall the definition (2.1) leading to the bounds,
\[
| \sum_{i_k} J_{i_1 \cdots i_p} \left( \prod_{l=1}^{l-1} x^{i_l} \right) \left( \prod_{h=l+1}^{p-1} \bar{x}^{i_h} \right) \prod_{l=1}^p y^{i_l} | \leq ||J||_\infty [1 + \left( \frac{1}{N} ||x||^2 \right)^{\frac{m-2}{2}}] [1 + \left( \frac{1}{N} ||\bar{x}||^2 \right)^{\frac{m-2}{2}}] ||x - \bar{x}|| ||y||,
\]
for \( 1 \leq l \leq p-1 \leq m-1 \). Consequently, we get for some \( c = c(a_1, \ldots, a_m) < \infty \) which is independent of \( N \), that
\[
\sum_{i=1}^N (G^i(x) - G^i(\bar{x})) y^i \leq c ||J||_\infty [1 + \left( \frac{1}{N} ||x||^2 \right)^{\frac{m-2}{2}}] [1 + \left( \frac{1}{N} ||\bar{x}||^2 \right)^{\frac{m-2}{2}}] ||x - \bar{x}|| ||y||,
\]
for all \( x, \bar{x} \) and \( y \). Taking \( y^i = G^i(x) - G^i(\bar{x}) \) results with the bound of (2.10), out of which we get the bound (2.11) by considering \( \bar{x} = 0 \).

By the estimate (B.6) of Appendix B we know that for any \( k < \infty \),
\[
\sup_N \mathbb{E} \left[ ||J||_\infty^k \right] < \infty,
\]
for the norm \( ||J||_\infty \) of (2.1). By the boundedness of \( N \mapsto \mathbb{E}[K_N(0)^k] \) and (2.2) the bound (2.12) immediately implies that for each \( k < \infty \), also
\[
\sup_N \mathbb{E} \left[ \sup_{t \in \mathbb{R}^+} K_N(t)^k \right] < \infty.
\]

Building upon (2.12) and (2.13) we next prove uniform moment bounds and pre-compactness of the functions of interest to us.
Proposition 2.3. Let $U_N$ denote any one of the functions $A_N$, $F_N$, $\chi_N$ and $C_N$. Then, $\mathbb{E}[\sup_{s,t \leq T}|U_N(s,t)|^k]$ is uniformly bounded in $N$, for any fixed $T < \infty$ and $k < \infty$. Further, for any fixed $T < \infty$, the sequence of continuous functions $U_N(s,t)$ is pre-compact almost surely and in expectation with respect to the uniform topology on $[0,T]^2$.

Proof: We start with the uniform bound on the moments of $A_N$, $F_N$, $\chi_N$ and $C_N$. To this end, let $B_N(t) := \frac{1}{N} \sum_{i=1}^N |B_i|^2$ and $G_N(t) := \frac{1}{N} \sum_{i=1}^N |G(x_i)|^2$. Fixing $T, k < \infty$ let $\|K_N\|_{\infty} := \sup\{K_N(t) : 0 \leq t \leq T\}$, and similarly define $\|B_N\|_{\infty}$, $\|G_N\|_{\infty}$ and $\|U_N\|_{\infty} := \mathbb{E}[\sup_{0 \leq s,t \leq T}|U_N(s,t)|]$.

A key estimate is then the following,

$$\mathbb{E}[\|B_N\|_{\infty}^k] + \mathbb{E}[\|K_N\|_{\infty}^k] + \mathbb{E}[\|B_N\|_{\infty}^k] + \mathbb{E}[\|G_N\|_{\infty}^k] < \infty,$$

holding for each fixed $k$. Indeed, the bounds on $\|B_N\|_{\infty}$ and $\|K_N\|_{\infty}$ are already obtained in (2.12) and (2.13), whereas by (2.11) we have that

$$G_N(t) \leq c(1 + K_N(t)^{\frac{k-1}{2}}),$$

yielding by (2.12) and (2.13) the uniform moment bound on $\|G_N\|_{\infty}$. Finally, it is easy to show that

$$\mathbb{P}(\|B_N\|_{\infty} \geq 4T(3 + x)) \leq e^{-Nx},$$

for all $T, x > 0$ (c.f. [3, (3.44)]), thereby providing a uniform bound for each moment of $\|B_N\|_{\infty}$ and deriving the conclusion of (2.14).

Similarly, by (2.3), (2.13), (2.10) and the exponential in $N$ bound of (2.17) on the tail of $\|J\|_{\infty}$, we have for each $L > 0$, that there exists $M = M(L) < \infty$ such that for all $N$,

$$\mathbb{P}(\|J\|_{\infty} + \|B_N\|_{\infty} + \|G_N\|_{\infty} \geq M) \leq e^{-Lx}.$$
for \( t, t' \in [i\delta, (i+2)\delta], i = 0, 1, \ldots, T/\delta \) and applying (2.10) with \( T = 2\delta \). In case \( b = G(x) \), note that by (2.10) we have that
\[
\frac{1}{N} \sum_{i=1}^{N} (G^i(x_i) - G^i(x_{i'}))^2 \leq 4c^2(\|J\|_\infty^2) (1 + K_N(t)(m-2))(1 + K_N(t')(m-2))
\]
for all \( t, t' \). Thus, in view of the bounds (2.11) and (2.17), everything reduces to showing that
\[
\mathbb{P} \left( \sup_{|t-t'|<\delta} \mathbb{E} \left[ \frac{1}{N} \sum_{i=1}^{N} |x_i^t - x_{i'}^t|^2 > \epsilon \right] \right) \leq e^{-L'(\delta, \epsilon)N}, \quad \mathbb{E} \left[ \left( \frac{1}{N} \sum_{i=1}^{N} |x_i^t - x_{i'}^t|^2 \right)^2 \right] \leq L'(\delta, \epsilon)^{-1}
\]
for some \( L' \) with the same properties as \( L \). To this end, note that by (1.6)
\[
|x_i^t - x_{i'}^t| \leq |B_i^t - B_{i'}^t| + \|f'(K_N)\| \int_t^{t'} |x_u^t| du + \int_t^{t'} |G^i(x_u)| du.
\]
So by (1.4), for some universal constant \( \rho_1 < \infty \), all \( t, t' \) and \( N \),
\[
\frac{1}{N} \sum_{i=1}^{N} |x_i^t - x_{i'}^t|^2 \leq \frac{3}{N} \sum_{i=1}^{N} |B_i^t - B_{i'}^t|^2 + 3|t - t'|^2 \left[ \rho_1 (1 + \|K_N\|_{\infty})^{2r} \|K_N\|_{\infty} + \|G_N\|_{\infty} \right]
\]
Thus, by (2.14), (2.16) and (2.17), we readily obtain (2.18). \( \square \)

A key ingredient in the derivation of the limiting dynamics is the “self-averaging” of the functions \( A_N \), \( F_N \), \( C_N \), \( \chi_N \) (and hence of \( D_N \) and \( E_N \)), as we next state and prove.

**Proposition 2.4.** Assume that the concentration of measure of Hypothesis 1.1 holds, and as before let \( U_N \) denote any one of the functions \( A_N \), \( F_N \), \( C_N \), \( \chi_N \). Then, for any \( T < \infty \) and \( \rho > 0 \),
\[
\sum_N \mathbb{P} \left[ \sup_{s,t \leq T} |U_N(s, t) - \mathbb{E}(U_N(s, t))| \geq \rho \right] < \infty,
\]
implies by the uniform moment bounds on \( \|U_N\|_{\infty} \) that
\[
\lim_{N \to \infty} \sup_{s,t \leq T} \mathbb{E} \left[ |U_N(s, t) - \mathbb{E}(U_N(s, t))|^2 \right] = 0.
\]

Our strategy relies on the following general ”localized” version of the concentration of measure property, which might be of some independent interest.

**Lemma 2.5.** Suppose \( V_N \) are functions on normed spaces \( \mathcal{E}_N \) in which the concentration of measure of Hypothesis 1.1 holds. Assume that \( K = \sup_{N} \mathbb{E}[V_N^2]^{1/2} < \infty \) and that for every \( M < \infty \) there exists a subset \( \mathcal{L}_{N,M} \) of \( \mathcal{E}_N \) on which \( |V_N| \leq 2M \) and \( V_N \) is Lipschitz with Lipschitz constant \( A_{N,M} \leq \frac{D(M)}{\sqrt{N}} \). Further assume that, for every \( L > 0 \) there exists an \( M = M(L) \) such that \( \mathbb{P}[\mathcal{L}_{N,M}] \leq \exp(-LN) \) for all \( N \). Then,
\[
\mathbb{P}[|V_N - \mathbb{E}[V_N]| \geq \rho] \leq C^{-1} \exp(-C(\frac{\rho}{2D(M(L))})^N) + 4(K + M(L))\rho^{-1}e^{-LN/2} + e^{-NL}.
\]

**Proof:** Recall that for every \( M < \infty \) there exists a globally Lipschitz function
\[
U^M_N(z) = \sup_{y \in \mathcal{L}_{N,M}} \{V_N(y) - A_{N,M}||z - y||\}
\]
on \( \mathcal{E}_N \), of Lipschitz constant \( A_{N,M} \), which coincides with \( V_N \) on the set \( \mathcal{L}_{N,M} \). These properties are inherited by \( V^M_N = \max(U^M_N, -2M) \) for which also \( |V^M_N| \leq 2M \). We thus get (2.21) by combining the triangle inequality
\[
|V_N - \mathbb{E}[V_N]| \leq |V_N - V^M_N| + |V^M_N - \mathbb{E}[V^M_N]| + |\mathbb{E}[V^M_N] - \mathbb{E}[V_N]|,
\]
with the Cauchy-Schwartz inequality

$$|E[V_N - V_N^M]| \leq E|V_N - V_N^M| \leq 2M P(\mathcal{L}_{N,M}^c) + E[|V_N|1_{\mathcal{L}_{N,M}^c}] \leq (K + 2M)P[\mathcal{L}_{N,M}^c]^{1/2},$$

and applying the concentration of measure inequality of Hypothesis 1.1 for $V_N^M$. □

**Proof of Proposition 2.4:** We wish to apply the estimate (2.21) to $V_N = U_N(s, t)$ for any of our four functions, and any fixed pair of times $s, t$. To this end, for each $M < \infty$ and any $N$ define the subset

$$\mathcal{L}_{N,M} = \{(x_0, J, B) \in \mathcal{E}_N : ||J||_{\infty} + \|B_N\|_{\infty} + \|K_N\|_{\infty} + \|G_N\|_{\infty} \leq M \}$$

of $\mathcal{E}_N$. For $M$ sufficiently large, the probability of the complement set $\mathcal{L}_{N,M}^c$ decays exponentially in $N$ by (2.17), whereas by Proposition 2.3 we have the uniform moment bounds for the functions $U_N(s, t)$, as well as the stated pointwise bound in $\mathcal{L}_{N,M}$. It thus suffices to prove the stated Lipschitz constant of $V_N$ on $\mathcal{L}_{N,M}$. We start by showing next that restricted to the set $\mathcal{L}_{N,M}$, the solution $x$ of (1.6) is a Lipschitz function of the triple $(x_0, J, B)$.

**Lemma 2.6.** Let $x, \tilde{x}$ be the two strong solutions of (1.6) constructed from $(x_0, J, B)$ and $(\tilde{x}_0, \tilde{J}, \tilde{B})$, respectively. If $(x_0, J, B)$ and $(\tilde{x}_0, \tilde{J}, \tilde{B})$ are both in $\mathcal{L}_{N,M}$, then we have the Lipschitz estimate:

$$\sup_{t \leq T} \frac{1}{N} \sum_{i \leq N} |x_i^t - \tilde{x}_i^t|^2 \leq \frac{D_0(M, T)}{N} \|(x_0, J, B) - (\tilde{x}_0, \tilde{J}, \tilde{B})\|^2,$$

where the finite constant $D_0(M, T)$ is independent of $N$.

**Proof:** We denote by $G^i(\cdot)$ (resp. $\tilde{G}^i(\cdot)$) the Gaussian fields constructed from the $J$ (resp. $\tilde{J}$). We write the following natural decomposition:

$$e_N(t) := \frac{1}{N} \sum_{i=1}^N |x_i^t - \tilde{x}_i^t|^2 = \frac{1}{N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)(x_i^t - \tilde{x}_i^t) - \int_0^t (f'(K_N(u)) - f'(\tilde{K}_N(u)))x_i^t \, du - \int_0^t f'(\tilde{K}_N(u))(x_i^t - \tilde{x}_i^t) \, du$$

$$+ \int_0^t (G^i(x_u) - G^i(\tilde{x}_u)) \, du + \int_0^t (G^i(\tilde{x}_u) - \tilde{G}^i(\tilde{x}_u)) \, du$$

$$= I_1 + I_2 + \ldots + I_6$$

Then, since $\|K_N\|_{\infty} \leq M$ and $\|\tilde{K}_N\|_{\infty} \leq M$, we have that for some finite $C(M, T)$ that is independent of $N$ and for all $c > 0$,

$$I_1 \leq \frac{c}{2N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2 + \frac{1}{2cN} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2$$

$$I_2 \leq \frac{c}{2N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2 + \frac{1}{2cN} \sum_{i=1}^N (B_i^t - \tilde{B}_i^t)^2$$

$$I_3 + I_4 \leq C(M, T) \left( \frac{c}{2N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2 + \int_0^t \frac{1}{2cN} \sum_{i=1}^N (x_u^t - \tilde{x}_u^t)^2 \, du \right)$$

$$I_6 \leq C(M, T) \left( \frac{c}{2N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2 + \int_0^t \frac{1}{2cN} \sum_{i=1}^m \sum_{p=1}^{p=1} \sum_{1 \leq n \leq t \leq N} (N^{p-1}\sum_{i=1}^p (J_{n+i} - \tilde{J}_{n+i})^2) \right)$$
We can bound the term $I_5$ on $L_{N,M}$, using (2.10) of Lemma 2.2 by

$$I_5 \leq C(M)\|J\|_\infty^N \int_0^t \left( \frac{1}{N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2 \right)^{\frac{1}{2}} \left( \frac{1}{N} \sum_{i=1}^N (\tilde{x}_i^u - \tilde{x}_i^t)^2 \right)^{\frac{1}{2}} du$$

$$\leq C(M,T) \left( \frac{c}{2N} \sum_{i=1}^N (x_i^t - \tilde{x}_i^t)^2 + \int_0^t \frac{1}{2cN} \sum_{i=1}^N (x_i^u - \tilde{x}_i^u)^2 du \right),$$

for some $C(M,T)$ independent of $N$ and all $c > 0$. Adding these estimates we get the bound

$$e_N(t) \leq \tilde{C}(M,T) \left[ c e_N(t) + \frac{1}{cN} \|(x_0, J, B) - (\tilde{x}_0, \tilde{J}, \tilde{B})\|^2 + \frac{1}{c} \int_0^t e_N(u) du \right],$$

on $e_N(t)$ of (2.23), so for $c = c(M,T) > 0$ small enough, Gronwall’s lemma yields the stated bound.

Equipped with Lemma 2.6 it is now easy to prove that

**Lemma 2.7.** Let $x, \tilde{x}$ be the two strong solutions of (1.4) constructed from $(x_0, J, B)$ and $(\tilde{x}_0, \tilde{J}, \tilde{B})$, respectively. If $(x_0, J, B)$ and $(\tilde{x}_0, \tilde{J}, \tilde{B})$ are both in $L_{N,M}$, then we have the Lipschitz estimate for each of the four functions $U_N(s,t)$ of interest,

$$(2.24) \quad \sup_{s,t \leq T} |U_N(s,t) - \tilde{U}_N(s,t)| \leq \frac{D(M,T)}{\sqrt{N}} \|(x_0, J, B) - (\tilde{x}_0, \tilde{J}, \tilde{B})\|,$$

where the constant $D(M,T)$ depends only on $M$ and $T$ and not on $N$.

**Proof:** Since each of the four functions $U_N(s,t)$ is of the form $\frac{1}{N} \sum_{i=1}^N a_s^i b_t^i$, we have that

$$|U_N(s,t) - \tilde{U}_N(s,t)| \leq \frac{1}{N} \sum_{i=1}^N |a_s^i - \tilde{a}_s^i||b_t^i| + \frac{1}{N} \sum_{i=1}^N |\tilde{a}_s^i||b_t^i - \tilde{b}_t^i|$$

$$\leq \left[ \frac{1}{N} \sum_{i=1}^N |a_s^i - \tilde{a}_s^i|^2 \right]^{1/2} \left[ \frac{1}{N} \sum_{i=1}^N |b_t^i|^2 \right]^{1/2} + \left[ \frac{1}{N} \sum_{i=1}^N |b_t^i - \tilde{b}_t^i|^2 \right]^{1/2} \left[ \frac{1}{N} \sum_{i=1}^N |\tilde{a}_s^i|^2 \right]^{1/2}.$$

Here the functions $a_t$ and $b_t$ are either $x_t$, $B_t$ or $G(x_t)$. This bound and Lemma 2.6 are sufficient to prove the conclusion of the lemma for the two functions $C_N(s,t)$ and $\chi_N(s,t)$. To prove it for the other two functions $A_N(s,t)$ and $F_N(s,t)$, note that by (2.10) we have the bound

$$\left[ \frac{1}{N} \sum_{i=1}^N |G^i(x_s) - G^i(\tilde{x}_s)|^2 \right]^{1/2} \leq C(M)\|J\|_\infty^N (1 + M^{\frac{p-1}{2}}) \left( \frac{1}{N} \sum_{1 \leq i \leq N} |x_s^i - \tilde{x}_s^i|^2 \right)^{1/2},$$

holding on $L_{N,M}$, so combining the obvious consequence of Cauchy-Schwartz

$$\left[ \frac{1}{N} \sum_{i=1}^N |G^i(\tilde{x}_s) - \tilde{G}^i(\tilde{x}_s)|^2 \right]^{1/2} \leq C(M)\left[ \frac{1}{N} \sum_{p=1}^m \sum_{1 \leq i_1 \cdots i_p \leq N} (N^{\frac{p-1}{2}} (J_{i_1 \cdots i_p} - \tilde{J}_{i_1 \cdots i_p})^2 \right]^{1/2},$$

on $L_{N,M}$, with Lemma 2.6 one gets

$$\left[ \frac{1}{N} \sum_{i=1}^N |G^i(x_s) - \tilde{G}^i(\tilde{x}_s)|^2 \right]^{1/2} \leq \frac{C(M,T)}{\sqrt{N}} \|(x_0, J, B) - (\tilde{x}_0, \tilde{J}, \tilde{B})\|.$$

Moreover, the estimate (2.11) gives the bound
\[
\left( \frac{1}{N} \sum_{i=1}^{N} |G^i(x_i)|^2 \right)^{1/2} \leq c||J||_\infty^N (1 + M^{m-1}) \leq C(M).
\]

The last two estimates and Lemma 2.6 thus yield the conclusion (2.21) for both \(A_N(s,t)\) and \(F_N(s,t)\). □

In view of Lemma 2.7, the inequality (2.21) applies for \(V_N = U_N(s,t)\), for any fixed \(s,t \leq T\) with constants \(K\) and \(D = D(M(L),T)\) that are independent of \(s,t,\rho\) and \(N\). Consequently, by the union bound, for any finite subset \(\mathcal{A}\) of \([0,T]^2\), and any \(\rho > 0\), the sequence \(N \to \mathbb{P}[\sup_{(s,t) \in \mathcal{A}} |U_N(s,t) - E U_N(s,t)| \geq \rho/3] \) is summable.

Recall that in the course of proving Proposition 2.3 we showed that for any \(\epsilon > 0\) there exists \(\bar{L}(\delta,\epsilon) \to \infty\) for \(\delta \to 0\), such that for all \(N\),
\[
\mathbb{P}( \sup_{|s-s'| + |t-t'| < \delta} |U_N(s,t) - U_N(s',t')| > \epsilon ) \leq e^{-\bar{L}(\delta,\epsilon)N}, \quad \sup_{|s-s'| + |t-t'| < \delta} |EU_N(s,t) - EU_N(s',t')| \leq \bar{L}(\delta,\epsilon)^{-1}.
\]

Choosing \(\delta > 0\) small enough so that \(\bar{L}(2\delta, \rho/3) > 3/\rho > 0\), we thus get (2.19) by considering the finite subset \(\mathcal{A} = \{(i\delta, j\delta) : i,j = 0,1,\ldots,T/\delta\}\) of all points of \([0,T]^2\) on a \(\delta\)-mesh. □

We shall often apply the following direct consequence of Propositions 2.3 and 2.4:

**Corollary 2.8.** Suppose that \(\Psi : \mathbb{R}^\ell \to \mathbb{R}\) is locally Lipschitz with \(|\Psi(z)| \leq M||z||_k^\ell\) for some \(M,\ell,k < \infty\), and \(Z_N \in \mathbb{R}^\ell\) is a random vector, where for \(j = 1,\ldots,\ell\), the \(j\)-th coordinate of \(Z_N\) is one of the functions \(A_N, F_N, \chi_N\) or \(C_N\) evaluated at some \((s_j,t_j) \in [0,T]^2\). Then,
\[
\lim_{N \to \infty} \sup_{s_j,t_j} |\mathbb{E}[\Psi(Z_N) - \Psi(E Z_N)]| = 0.
\]

**Proof:** It follows from Proposition 2.3 that \(R = \sup_{s_j,t_j,N} ||E(Z_N)||_k < \infty\). For each \(r \geq R\) let \(c_r\) denote the finite Lipschitz constant of \(\Psi(\cdot)\) (with respect to \(||\cdot||_2\) on the compact set \(\Gamma_r := \{z : ||z||_k \leq r\}\). Then,
\[
|\mathbb{E}[\Psi(Z_N) - \Psi(E Z_N)]| \leq \mathbb{E}[|\Psi(Z_N) - \Psi(E Z_N)|1_{Z_N \in \Gamma_r} + |\Psi(Z_N)|1_{Z_N \notin \Gamma_r}] \leq c_r E(||Z_N - E Z_N||_2) + 2\ell M r^{-k} E(||Z_N||_k^\ell).\]

We have by (2.20) and the uniform moment bounds of Proposition 2.3 that \(\sup_{s_j,t_j} E(||Z_N - E Z_N||_2) \to 0\) as \(N \to \infty\), while \(c' = \sup_{s_j,t_j,N} E||Z_N||_k^\ell < \infty\), implying that
\[
\lim_{N \to \infty} \sup_{s_j,t_j} |\mathbb{E}[\Psi(Z_N) - \Psi(E Z_N)]| \leq 2c' \ell M r^{-k},
\]
which we make arbitrarily small by taking \(r \to \infty\). □

**3. LIMITING EQUATIONS: PROOF OF PROPOSITION 1.3**

We shall denote in short
\[
C_N^\alpha(s,t) = \mathbb{E}[C_N(s,t)], \quad \chi_N^\alpha(s,t) = \mathbb{E}[\chi_N(s,t)]
\]
where expectation is over the Brownian path \(B\), the disorder \(J\) and the initial condition \(x_0\), and adopt a similar notation for other functions of interest.

Integrating the SDS (1.6) we have that
\[
x^t_s = x^t_0 + B^t_s - \int_0^s f'(K_N(u))x^t_u du + \int_0^s G^t(x_u)du.
\]
Hence, upon multiplying by \(x_i^t\) and \(B_i^t\) followed by averaging over \(i\) and taking the expected value, we get that for any \(s, t \in \mathbb{R}^+\),

\[
C_N^a(s, t) = C_N^a(0, t) + \chi_N^a(t, s) - \int_0^s \mathbb{E}[f'(K_N(u))C_N(u, t)]du + \int_0^s A_N^a(u, t)du
\]

(3.2)

\[
\chi_N^a(s, t) = \chi_N^a(0, t) + t \wedge s - \int_0^s \mathbb{E}[f'(K_N(u))\chi_N(u, t)]du + \int_0^s F_N^a(u, t)du.
\]

(3.3)

In the following, we use \(a_N \approx b_N\) when \(a_N(\cdot, \cdot) - b_N(\cdot, \cdot) \to 0\) as \(N \to \infty\), uniformly on \([0, T]^2\). Applying Corollary 2.8 for \(\Psi(z) = z_1 f'(z_2)\) whose polynomial growth is guaranteed by our assumption (1.3), we deduce that

\[
\mathbb{E}[f'(K_N(u))C_N(u, t)] \approx f'(K_N^a(u))C_N^a(u, t), \quad \mathbb{E}[f'(K_N(u))\chi_N(u, t)] \approx f'(K_N^a(u))\chi_N^a(u, t).
\]

Our next proposition, approximates the terms \(A_N^a\) and \(F_N^a\) which we need in order to compute the limits of (3.2) and (3.3) as \(N \to \infty\).

**Proposition 3.1.** We have that

\[
A_N^a(t, s) \approx \nu'(C_N^a(t, t \vee s)) C_N^a(s, t \wedge s) - \nu'(C_N^a(t, 0))C_N^a(s, 0) - \int_0^{s \vee t} \nu'(C_N^a(t, u))D_N^a(s, u)du.
\]

(3.4)

\[
- \int_0^{s \vee t} \nu''(C_N^a(t, u))C_N^a(s, u)D_N^a(u, t)du.
\]

Further,

\[
F_N^a(t, s) \approx \chi_N^a(s, t \wedge s)\nu'(C_N^a(s, s)) - \int_0^{t \wedge s} \nu'(C_N^a(s, u))du - \int_0^s \nu'(C_N^a(s, u))E_N^a(u, t \wedge u)du
\]

(3.5)

\[- \int_0^s \chi_N^a(u, t \wedge u)\nu''(C_N^a(s, u))D_N^a(u, s)du.
\]

Deferring the proof of Proposition 3.1 we first use its conclusion to complete the proof of Proposition 1.3. To this end, note that \(B_0 = 0\) so \(F_N^a(s, 0) = 0\), for all \(N\) and \(s\). Further, with \(G_s = \sigma(J, x_0, B_u, u \leq s)\), we have by (1.15) that for all \(N\) and \(t > s\),

\[
F_N^a(s, t) = F_N^a(s, s) + \frac{1}{N} \sum_{i=1}^N \mathbb{E}[G^i(x_s)\mathbb{E}(B_i^j - B_i^j|G_s)] = F_N^a(s, s)
\]

(recall the independence of \(B_t - B_s\) and \(G_s\)). By the same reasoning also \(\chi_N^a(s, 0) = 0, E_N^a(s, 0) = 0, E_N^a(s, t) = E_N^a(s, s)\) and \(\chi_N^a(s, t) = \chi_N^a(s, s)\) for all \(N\) and any \(t \geq s \geq 0\) (cf. (1.9) and (1.16)). Likewise, by definition \(C_N^a(t, s) = C_N^a(s, t)\).

Let \(\Phi(C, \chi, D, E) : \mathcal{C}([0, T]^2)^4 \to \mathcal{C}([0, T]^2)^4\) denote the difference between the two sides of equations (1.17) - (1.20). In view of the above boundary and symmetry conditions, comparing (3.2) - (3.5) with (1.17) - (1.20) we see that \(\Phi(C_N^a, \chi_N^a, D_N^a, E_N^a) \to 0\) as \(N \to \infty\), uniformly on \([0, T]^2\). It is not hard to check that \(\Phi(\cdot)\) is continuous with respect to the topology of uniform convergence, hence \(\Phi(\cdot) = 0\) at any limit point of \((C_N^a, \chi_N^a, D_N^a, E_N^a)\), which also necessarily satisfies the same boundary and symmetry conditions, thus completing the proof of Proposition 1.3. □

As already noted, the first step in proving Proposition 3.1 is,

**Lemma 3.2.** Let \(\mathbb{E}_J\) denotes the expectation with respect to the Gaussian law \(\mathbb{P}_J\) of the disorder \(J\). Then, for each continuous path \(x \in \mathcal{C}(\mathbb{R}_+, \mathbb{R}^N)\) and all \(s, t \in [0, T]\) and \(i, j \in \{1, \ldots, N\}\),

\[
k_{ij}^s(t) := \mathbb{E}_J[G^i(x_t)G^j(x_s)] = \frac{x_i^t x_j^s}{N} \nu''(C_N(s, t)) + 1_{i=j} \nu'(C_N(s, t)).
\]

(3.6)
Proof: Observe that $G^i(x)$ of (3.10) are, for any given $x$, centered, jointly Gaussian random variables. Further, by our choice of $c\{i_1, \ldots, i_p\}$ it is not hard to verify that for any $i, i_1, \ldots, i_{p-1} = 1, \ldots, N$

(3.7) \[ c\{i, i_1, \ldots, i_{p-1}\} \sum_{1 \leq j_1, \ldots, j_{p-1} \leq N} 1_{\{i, i_1, \ldots, i_{p-1}\} = \{j_1, j_2, \ldots, j_{p-1}\}} = (p - 1)! (\sum_{r=1}^{p-1} 1_{j = j_r}). \]

Hence, by (3.7) and (3.8), we have that for any given vectors $x$ and $y$,

\[ \mathbb{E}[G^i(x)G^j(y)] = \sum_{p=1}^{m} \frac{a_p^2}{(p-1)!} \sum_{i_1, \ldots, i_{p-1}} \mathbb{E}(J_{i_1 \ldots i_{p-1}}J_{i_{p-1} \ldots i_{p-1}})x^{i_1} \ldots x^{i_{p-1}} y^{j_1} \ldots y^{j_{p-1}} \]

Further, by our choice of $(1.7)$ and $(3.7)$, we have that for any given vectors $x$ and $y$,

\[ \mathbb{E}[G^i(x)G^j(y)] = \sum_{p=1}^{m} \frac{a_p^2}{(p-1)!} N^{-(p-1)} \sum_{i_1, \ldots, i_{p-1}} (1 + \sum_{r=1}^{p-1} 1_{i = i_r}) x^{i_1} \ldots x^{i_{p-1}} y^{j_1} \ldots y^{j_{p-1}} = \sum_{p=1}^{m} \frac{a_p^2}{(p-1)!} \left( \frac{1}{N} \sum_{\ell=1}^{N} \frac{N}{p} \right)^{p-1} + \frac{x^{i_1}y^{j_1}}{N} \sum_{p=2}^{m} \frac{a_p^2}{(p-2)!} \left( \frac{1}{N} \sum_{\ell=1}^{N} \frac{N}{p} \right)^{p-2}, \]

so that with $\nu(\rho) = \sum_{p=1}^{m} \frac{a_p^2}{p!} \rho^p$ we have,

(3.8) \[ \mathbb{E}[G^i(x)G^j(y)] = \nu'\left( \frac{1}{N} \sum_{\ell=1}^{N} x^{j_1} \right) + \nu''\left( \frac{1}{N} \sum_{\ell=1}^{N} x^{j_1} \right). \]

Further, if $i \neq j$, then by (3.7) we also have that

\[ x^{i_1}y^{j_1} \mathbb{E}[G^i(x)G^j(y)] = \sum_{p=2}^{m} \frac{a_p^2}{(p-1)!} \sum_{i_1, \ldots, i_{p-1}} \mathbb{E}(J_{i_1 \ldots i_{p-1}}J_{i_{p-1} \ldots i_{p-1}})x^{i_1} \ldots x^{i_{p-1}} y^{i_1} \ldots y^{i_{p-1}} \]

Further, if $i \neq j$, then by (3.7) we also have that

\[ x^{i_1}y^{j_1} \mathbb{E}[G^i(x)G^j(y)] = \sum_{p=2}^{m} \frac{a_p^2}{(p-1)!} \sum_{i_1, \ldots, i_{p-1}} \mathbb{E}(J_{i_1 \ldots i_{p-1}}J_{i_{p-1} \ldots i_{p-1}})x^{i_1} \ldots x^{i_{p-1}} y^{i_1} \ldots y^{i_{p-1}} \]

implying that when $i \neq j$,

(3.9) \[ \mathbb{E}[G^i(x)G^j(y)] = \frac{x^{i_1}y^{j_1}}{N} \nu''\left( \frac{1}{N} \sum_{\ell=1}^{N} x^{j_1} \right), \]

so replacing $x$ and $y$ by $x_i$ and $x_s$ respectively, we immediately get (3.6) out of (3.8), (3.9) and the definition of $C_N(\cdot, \cdot)$.

Proof of Proposition 5.1 Fixing a continuous path $x$, let $k_t$ denote the operator on $L_2(\{1, \cdots N\} \times [0, t])$ with the kernel $k = k(x)$ of (3.10). That is, for $f \in L_2(\{1, \cdots N\} \times [0, t])$, $u \leq t$, $i \in \{1, \cdots, N\}$

(3.10) \[ [k_tf]_u^i = \sum_{j=1}^{N} \int_{0}^{t} s_{ij}^{2} f_j dv, \]
which is clearly also in $L_2(\{1, \cdots N\} \times [0, t])$. We next extend the definition \ref{eq:3.10} to the stochastic integrals of the form

$$[k_t \circ dZ]^i_u = \sum_{j=1}^N \int_0^t k^{ij}_{uv} dZ^j_v,$$

where $Z^j_v$ is a continuous semi-martingale with respect to the filtration $\mathcal{F}_t = \sigma(x_u : 0 \leq u \leq t)$ and is composed for each $j$, of a squared-integrable continuous martingale and a continuous, adapted, squared-integrable finite variation part. In doing so, recall that by \ref{eq:3.10} each $k^{ij}_{uv}(x)$ is the finite sum of terms such as $x^i_u \cdots x^j_v x^i_v \cdots x^j_v$, where in each term $a, b$ and $i_1, \ldots, i_a, j_1, \ldots, j_b$ are some non-random integers. Keeping for simplicity the implicit notation $\sum_{k=1}^\infty$ we thus adopt hereafter the convention of accordingly decomposing such integral to a finite sum, taking for each of its terms such as $\sum_{k=1}^\infty x^i_u \cdots x^j_v$ outside the integral, resulting with the usual Itô adapted stochastic integrals.

The latter are well defined, with $[k_t \circ dZ]^i_u$ being in $L_2(\{1, \cdots N\} \times [0, t])$ (recall Proposition \ref{prop:2.1} that $x_t$ has uniformly bounded finite moments of all orders under the joint law $\mathbb{P}_J \otimes \mathbb{P}^N_{x_0,J}$, hence so does the kernel $k^{ij}_{uv}(x)$).

Equipped with these definitions, we next claim that

\begin{equation}
\tag{3.11}
U^i_s(x) = x^i_s - x^i_0 + \int_0^s f'(K_N(u)) x^i_u du
\end{equation}

\begin{equation}
\tag{3.12}
Z^i_s(x) = U^i_s(x) - \int_0^s V^i_u(x) du,
\end{equation}

are both continuous semi-martingales with respect to the filtration $\mathcal{F}_s$, composed of squared-integrable continuous martingales and finite variation parts. Moreover, such choice satisfies for any $i$ and $s \in [0, \tau]$,

\begin{equation}
\tag{3.13}
V^i_s + [k_t \circ dU]^i_s = [k_t \circ dU]^i_s,
\end{equation}

and $V^i_s = [k_t \circ dZ]^i_s$ for any $i$ and all $s \leq \tau$. Further, for any $u, v \in [0, \tau]$ and $i, j \leq N$, let

\begin{equation}
\tag{3.14}
\Gamma^i_{uv}(x) := \mathbb{E} \left[ (G^i(x_u) - V^i_u(x))(G^j(x_v) - V^j_v(x)) \mid \mathcal{F}_\tau \right]
\end{equation}

Then, we can choose a version of $\Gamma^i_{uv}$ such that for any $s, v \leq \tau$ and all $i, l \leq N$,

\begin{equation}
\tag{3.15}
\sum_{j=1}^N \int_0^\tau k^{ij}_{su} \Gamma_{uv}^{jl} du + \Gamma^i_{uv} = k^{il}_{sv}.
\end{equation}

\textbf{Proof:} Since $U^i_s(x) = \int_0^s G^i(x_u) du + B^i_s$ (see \ref{eq:4.1}), the relation \ref{eq:3.12} between $\mathbb{E}[B^i_s \mid \mathcal{F}_t], U^i_s(x)$ and $\mathbb{E}[G^i(x_u) \mid \mathcal{F}_t]$ follows, as well as the continuity and integrability properties of the semi-martingales $U_s$ and $Z_s$. Using hereafter $G^i_s$ to denote $G^i(x_s)$, let

\begin{equation}
\tag{3.16}
L^N_s = \exp \left\{ \sum_{i=1}^N \int_0^\tau G^i_s dU^i_s(x) - \frac{1}{2} \sum_{i=1}^N \int_0^\tau (G^i_s)^2 ds \right\}.
\end{equation}

By Girsanov formula we have that the restriction to $\mathcal{F}_\tau$ satisfies,

$$\mathbb{P}^N_{x_0,J} \mid \mathcal{F}_\tau = L^N \mathbb{P}^N_{x_0,0} \mid \mathcal{F}_\tau,$$
Hence, with $\tau \geq s$, for any bounded $\mathcal{F}_\tau$-measurable random variable $\Phi$,

$$
(3.17) \quad E[G_s^i \Phi] = E_J E_{\mathcal{F}_{\tau \land s}}[G_s^i \Phi] = E_{\mathcal{F}_{\tau \land s}}[E_J[G_s^i L_{\tau}^N] \Phi] = E\left[\frac{E_J[G_s^i L_{\tau}^N]}{E_J[L_{\tau}^N]} \Phi\right],
$$

where the right-most identity is due to the change of measure formula $Q_{x_0}^N = E_J(L_{\tau}^N)E_{\mathcal{F}_{\tau \land s}}$, restricted to $\mathcal{F}_\tau$. With $(3.17)$ holding for all bounded $\mathcal{F}_\tau$-measurable $\Phi$, it follows that,

$$
V_s^i = E[G_s^i | \mathcal{F}_\tau] = \frac{E_J[G_s^i L_{\tau}^N]}{E_J[L_{\tau}^N]},
$$

and the identity $(3.13)$ follows by the Gaussian change of measure identity $(C.2)$ of Proposition C.1. Noting that $k$ (suppressing the dependence of $E$)

$$
\text{Hence, with } \tau \geq s, \text{ for any bounded } \mathcal{F}_\tau \text{-measurable random variable } \Phi,
$$

$$
(3.17) \quad E[G_s^i \Phi] = E_J E_{\mathcal{F}_{\tau \land s}}[G_s^i \Phi] = E_{\mathcal{F}_{\tau \land s}}[E_J[G_s^i L_{\tau}^N] \Phi] = E\left[\frac{E_J[G_s^i L_{\tau}^N]}{E_J[L_{\tau}^N]} \Phi\right],
$$

where the right-most identity is due to the change of measure formula $Q_{x_0}^N = E_J(L_{\tau}^N)E_{\mathcal{F}_{\tau \land s}}$, restricted to $\mathcal{F}_\tau$. With $(3.17)$ holding for all bounded $\mathcal{F}_\tau$-measurable $\Phi$, it follows that,

$$
V_s^i = E[G_s^i | \mathcal{F}_\tau] = \frac{E_J[G_s^i L_{\tau}^N]}{E_J[L_{\tau}^N]},
$$

and the identity $(3.13)$ follows by the Gaussian change of measure identity $(C.2)$ of Proposition C.1. Noting that $dZ = dU - V$ (see $(3.12)$), we have by $(3.13)$ that for all $i$ and $s \in [0, \tau]$,

$$
[k_\tau \circ dZ]^i_s = [k_\tau \circ dU]^i_s - [k_\tau V]^i_s = V_s^i
$$

as claimed.

We now apply $(3.13)$ to derive $(3.4)$, the easy part of Proposition 3.1. To this end, fixing $s, t \in [0, T]^2$ let

$$
\hat{A}_N(t, s) = \frac{1}{N} \sum_{i=1}^{N} V_i^i(x) x_s^i,
$$

for $\tau = t \lor s$, noting that since $x_s^i$ is measurable on $\mathcal{F}_\tau$, by $(1.15)$,

$$
A_s^i(t, s) = E\left[\frac{1}{N} \sum_{i=1}^{N} E[G_s^i x_s^i | \mathcal{F}_\tau]\right] = E[\hat{A}_N(t, s)] = \hat{A}_N^s(t, s).
$$

With $t \leq \tau$, combining $(3.13)$ and $(3.11)$ we get that

$$
\hat{A}_N(t, s) + \frac{1}{N} \sum_{i,j=1}^{N} \int_0^T x_s^i k_{iu}^{ij} V_u^j du = \frac{1}{N} \sum_{i,j=1}^{N} \int_0^T f'(K_N(u)) x_s^i k_{iu}^{ij} x_u^j du + \frac{1}{N} \sum_{i,j=1}^{N} \int_0^T x_s^i k_{iu}^{ij} dx_u
$$

(suppressing the dependence of $k_{iu}^{ij}$ and $V_u^j$ on $x$, and following our convention regarding stochastic integrals such as $\int_0^T x_s^i k_{iu}^{ij} dx_u$). Using the explicit expression of $k_{iu}^{ij}(x)$, and collecting terms while changing the order of summation and integration, we arrive at the identity,

$$
\hat{A}_N(t, s) + \int_0^T C_N(s, u) \nu''(C_N(t, u)) \hat{A}_N(u, t) du + \int_0^T \nu'(C_N(t, u)) \hat{A}_N(u, s) du
$$

$$
= \int_0^T f'(K_N(u)) C_N(s, u) \nu''(C_N(t, u)) C_N(u, t) du + \int_0^T f'(K_N(u)) \nu'(C_N(t, u)) C_N(u, s) du
$$

$$
(3.18) + \int_0^T C_N(s, u) \nu''(C_N(t, u)) d_u C_N(u, t) + \int_0^T \nu'(C_N(t, u)) d_u C_N(u, s).
$$
Applying Lemma \[A.1\] for the semi-martingales $x = y = z = w = x$, and the polynomials $Q(\rho) = \nu'(\rho)$ and $P(\rho) = \rho$ evaluated at $\sigma = \tau$, $\theta = s$ and $v = t$, we replace the stochastic integrals of (3.18) by

$$
\nu'(C_N(\tau, t))C_N(\tau, s) - \nu'(C_N(0, t))C_N(0, s)
$$

(3.19)

$$
- \frac{1}{2N}C_N(t, t) \int_0^t \nu''(C_N(u, t))du - \frac{1}{N}C_N(s, t) \int_0^\tau \nu'(C_N(u, t))du.
$$

Clearly, $\hat{A}_N(t, s) = \mathbb{E}[A_N(t, s)|\mathcal{F}_\tau]$ has the same uniform moment bounds of Proposition \[2.3\] as $A_N$ and further inherits the self-averaging property \[2.20\] from $A_N$. Hence, we may and shall apply Corollary \[2.8\] with possibly $\hat{A}_N$ as one of the arguments of the locally Lipschitz function $\Psi(z)$ of at most polynomial growth at infinity. Doing so for the functions $z_1z_2\nu''(z_3)$, and $z_1\nu'(z_2)$ and applying Corollary \[2.8\] also for $f'(z_1)z_2\nu''(z_3)$ and $f'(z_1)\nu'(z_2)z_3$, upon utilizing the uniform convergence with respect to the points $(s_j, t_j) \subseteq [0, T]^2$, we deduce from \[3.18\] and \[3.19\] that

$$
\hat{A}_N^a(t, s) + \int_0^\tau C_N^a(s, u)\nu''(C_N^a(t, u))\hat{A}_N^a(u, t)du + \int_0^\tau \nu'(C_N^a(t, u))\hat{A}_N^a(u, s)du
$$

$$
\simeq \int_0^\tau f'(K_N^a(u))C_N^a(s, u)\nu''(C_N^a(t, u))C_N^a(u, t)du
$$

$$
+ \int_0^\tau f'(K_N^a(u))\nu'(C_N^a(t, u))C_N^a(s, u)du + \nu'(C_N^a(\tau, t))(C_N^a(\tau, s) - \nu'(C_N^a(0, t))C_N^a(0, s).
$$

Finally, recall that

$$
\hat{A}_N^a(t, s) = A_N^a(t, s) = D_N^a(s, t) + f'(K_N^a(t))C_N^a(s, t),
$$

with the corresponding replacement for $\hat{A}_N^a(u, t)$ and $\hat{A}_N^a(u, s)$. With $\tau = t \vee s$, we indeed arrive at \[3.11\].

We now turn to the more involved part of Proposition \[3.11\] namely, the derivation of \[3.5\]. To this end, as we have seen already, it suffices to consider $s \geq t$, as we do hereafter. To this end, taking $\tau = s$ we use the notation $V_u = \mathbb{E}[G_u|\mathcal{F}_s]$ of Lemma \[3.3\] while suppressing the dependence on $x$. Since $B_t = U_t - \int_0^t G_u dv$ with $U_t$ being $\mathcal{F}_s$-measurable, we deduce that

$$
\mathbb{E}[G^a_s B^i_t] = \mathbb{E}[G^a_s U^i_t] - \int_0^t \mathbb{E}[G^a_u G^a_s]dv = \mathbb{E}\left[\mathbb{E}[G^a_t|\mathcal{F}_s]U^i_t - \int_0^t V^i_u dv - \int_0^t \Gamma^i_{stu} dv\right],
$$

(recall that $\mathbb{E}(G_t - V_t|\mathcal{F}_s) = 0$ for all $v \leq s$ and $\Gamma^i_{stu} := \mathbb{E}[(G^i_u - V^i_u)(G^i_t - V^i_t)|\mathcal{F}_s]$ is per \[3.14\]). Further, by \[3.12\] and the identity $\mathbb{E}[G^a_t|\mathcal{F}_s] = [k_s \circ dZ^a_s]^t_s$ of Lemma \[3.3\] we get

$$
\mathbb{E}[G^a_s B^i_t] + \mathbb{E}\left[\int_0^t \Gamma^i_{stu} dv\right] = \mathbb{E}([k_s \circ dZ^a_s]^t_s Z^i_t).
$$

Since $Z^i_t = \mathbb{E}[B^i_t|\mathcal{F}_s]$ and $[k_s \circ dZ^a_s]$ is $\mathcal{F}_s$-measurable, we have that

$$
\mathbb{E}([k_s \circ dZ^a_s] Z^i_t) = \mathbb{E}((k_s \circ dZ^a_s) B^i_t) = \mathbb{E}(\mathbb{E}([k_s \circ dB]^a_s B^i_t|\mathcal{F}_s) B^i_t),
$$

where the right-most identity holds since the kernel of the linear operator $k_s$ is $\mathcal{F}_s$-measurable and $\mathbb{E}[B_u|\mathcal{F}_s] = Z_u$ for all $u \leq s$ (recall that $[k_s \circ dZ^a_s]$ is the $L_2$-limit of discrete sums with mash size
going to zero). In view of \(3.14\) and \(3.20\), we have that
\[
[k_s \circ dB_x]_t = [k_s \circ dx]_t + [k_s f'(K_N) x]_t - [k_s G]_t.
\]

(3.22)
\[
\int_0^s \nu'(C_N(s, u))dx_u + \int_0^s \nu''(C_N(s, u))x_u^i du + \int_0^s \nu''(C_N(s, u))f'(K_N(u))C_N(s, u)x_u^i du
- \int_0^s \nu''(C_N(s, u))x_u^i \Lambda_N(u, s) du - \int_0^s \nu'(C_N(s, u))G_u^i du
\]

Using Itô’s formula for \(x_u^i \nu'(C_N(s, u))\) we replace the two stochastic integrals in (3.22) by
(3.23)
\[
x_u^i \nu'(C_N(s, u)) - x_0^i \nu'(C_N(s, u)) - \frac{1}{2N} C_N(s, s) \int_0^s \nu''(C_N(s, u))x_u^i du - \frac{1}{N} \int_0^s \nu''(C_N(s, u)) du.
\]

Recall that by \(1.15\) and \(3.21\)
\[
F_N^t(s, t) + \mathbb{E}[\frac{1}{N} \sum_{i=1}^N \int_0^t \Gamma_{ss}^{ii} dv] = \mathbb{E}(\frac{1}{N} \sum_{i=1}^N \mathbb{E}([k_s \circ dB_x]_t | \mathcal{F}_s) B_t^i),
\]

which by the preceding is the expectation of
\[
\chi_N(s, t) \nu'(C_N(s, s)) - \chi_N(0, t) \nu'(C_N(s, 0))
- \frac{1}{2N} C_N(s, s) \int_0^s \nu''(C_N(s, u)) \chi_N(s, t) du - \frac{1}{N} \int_0^s \nu''(C_N(s, u)) \chi_N(s, t) du
+ \int_0^s f'(K_N(u)) \nu'(C_N(s, u)) \chi_N(0, t) du + \int_0^s f'(K_N(u)) \nu''(C_N(s, u)) C_N(s, u) \chi_N(s, t) du
\]

(3.24)
\[
- \int_0^s \nu''(C_N(s, u)) \chi_N(s, t) \Lambda_N(u, s) du - \int_0^s \nu'(C_N(s, u)) F_N(u, s) du + \kappa_N(s, t),
\]

where in view of \(1.15\)
\[
\kappa_N(s, t) := \frac{1}{N} \sum_{i=1}^N \int_0^s \nu'(C_N(s, u))(G_u^i - V_u^i) B_t^i du.
\]

Recall \(3.12\) that \(B_t^i = Z_t^i - \int_0^t (G_u^i - V_u^i) du\), hence
\[
\kappa_N(s, t) = \frac{1}{N} \sum_{i=1}^N \int_0^s \nu'(C_N(s, u))(G_u^i - V_u^i) Z_t^i du - \frac{1}{N} \sum_{i=1}^N \int_0^s \nu'(C_N(s, u)) \int_0^t (G_u^i - V_u^i)(G_v^i - V_v^i) dvdu.
\]

As both \(\nu'(C_N(u, s))\) and \(Z_t^i\) are \(\mathcal{F}_s\)-measurable while \(\mathbb{E}(G_u^i - V_u^i | \mathcal{F}_s) = 0\) for all \(u \leq s\), the expectation of the first term on the right-side vanishes. Further, conditioning on \(\mathcal{F}_s\), we have in view of \(3.14\) that
\[
\mathbb{E}[\kappa_N(s, t)] = - \mathbb{E}\left[ \frac{1}{N} \sum_{i=1}^N \int_0^s \nu'(C_N(s, u)) \int_0^t \Gamma_{uu}^{ii} du dv \right].
\]

All terms of \(3.24\) apart from \(\kappa_N\) are of the form covered by Corollary \(2.8\) for functions \(\Psi(z)\) similar to those encountered in the derivation of \(3.24\), namely, \(z_1 \nu'(z_2), z_1 \nu''(z_2), z_1 \nu'''(z_2), f'(z_1) \nu'(z_2), z_1 z_2 \nu''(z_3), f'(z_1) z_2 \nu''(z_3)\). Utilizing their uniform convergence and \(3.20\), while recalling that \(\chi_N(0, t) = 0\), \(3.20\) and the analogous
\[
F_N^t(u, t) = E_N^t(u, t) + f'(K_N(u)) C_N^a(s, u),
\]


is not hard to check that we get (3.5), once we prove the following lemma.

**Lemma 3.4.** For \( v \in [0, s] \), let

\[
\phi_N(s,v) := \nu'(C_N(s,v)) - \frac{1}{N} \sum_{i=1}^{N} \Gamma_{sv}^{ii} - \frac{1}{N} \sum_{i=1}^{N} \int_{0}^{s} \nu'(C_N(s,u)) \Gamma_{uv}^{ii} \, du
\]

(where \( \Gamma \) is defined with \( \tau = s \)). Then, for any \( t \leq s \),

\[
\int_{0}^{t} \mathbb{E}[\phi_N(s,v)] \, dv \simeq 0
\]

**Proof:** Fixing \( v \leq s \), recall that by (3.6) we have that

\[
\phi_N(s,v) + \frac{1}{N} \nu''(C_N(s,v)) C_N(s,v) = \frac{1}{N^2} \sum_{i,j=1}^{N} \int_{0}^{s} \nu''(C_N(s,u)) x_s^i x_u^j \Gamma_{uv}^{ij} \, du\]

Further, since we had \( \tau = s \), the right-side is identically zero by (3.15), implying by the \( \mathcal{F}_s \)-measurability of \( x_s, x_u \) and some algebraic manipulations that

\[
\phi_N(s,v) = \int_{0}^{s} \nu''(C_N(s,u)) \frac{1}{N^2} \sum_{i,j=1}^{N} x_s^i x_u^j \mathbb{E}[(G_i^u - V_i^u)(G_j^s - V_j^s) | \mathcal{F}_s] \, du\]

Consequently, with \( C_N(s,u) \) measurable on \( \mathcal{F}_s \) we have that

\[
\mathbb{E}[\phi_N(s,v)] = \int_{0}^{s} \mathbb{E}[\nu''(C_N(s,u)) (A_N(u,s) - \tilde{A}_N(u,s)) (A_N(v,u) - \tilde{A}_N(v,u)) | \mathcal{F}_s] \, du - \frac{1}{N} \mathbb{E}[\nu''(C_N(s,v)) C_N(s,v)].
\]

In view of Proposition 2.3 and (2.20), the first term converges to zero uniformly in \( s, \nu \) by the uniform \( L_2 \) convergence of \( A_N - A_N^a \) (hence also of \( \tilde{A}_N \)), and the uniform (in \( [0,T]^2 \) and \( N \)) bound on each moment of \( C_N \) and \( A_N \) (hence on those of \( \tilde{A}_N \) as well). By same reasoning, the second term converges to zero at rate \( 1/N \), uniformly in \( s, v \). Utilizing the uniformity of the convergence, we see that \( \int_{0}^{t} \mathbb{E}[\phi_N(s,v)] \, dv \simeq 0 \) as claimed. \( \square \)

### 4. Differentiability and uniqueness for the limiting dynamics

We start the proof of Theorem 1.2 by the next lemma relating the solutions of (1.17)-(1.20) with those of (1.12)-(1.14).

**Lemma 4.1.** Fixing \( T < \infty \), suppose \( (C, \chi, D, E) \) is a solution of the integral equations (1.17)-(1.20) in the space of bounded continuous functions on \([0,T]^2\) subject to the symmetry condition \( C(s,t) = C(t,s) \) and the boundary conditions \( E(s,0) = 0 \) for all \( s \), and \( E(s,t) = E(s,s) \) for all \( t \geq s \). Then, \( \chi(s,t) = \int_{0}^{t} R(s,u) \, du \) where \( R(s,t) = 0 \) for \( t > s \), \( R(s,s) = 1 \) and for \( T \geq s > t \), the bounded and absolutely continuous functions \( C, R \) and \( K(s) = C(s,s) \) necessarily satisfy the integro-differential equations (1.12)-(1.14).

**Proof:** Consider the integral operator \( k_C \) on \( C([0,T]) \) given by,

\[
[k_C h](s) := - \int_{0}^{s} \nu'(C(s,u)) h(u) \, du,
\]
and let
\[ h(s, t) := -f'(C(s, s))\chi(s, t) - \int_0^s \chi(u, t)\nu''(C(s, u))D(s, u)du + \chi(s, t)\nu'(C(s, s)) - \int_0^{t \wedge s} \nu'(C(s, u))du \]
Then, per fixed \( t \), the equation \( \text{(1.20)} \) states that \( E(s, t) = [k_C E(\cdot, t)](s) + h(s, t) \). Since the (continuous) kernel \( \nu'(C(s, u)) \) of \( k_C \) is uniformly bounded on \([0, T]^2\), it follows by Picard iterations (splitting \([0, T]\) to sufficiently small time intervals to guarantee convergence of the series \( \sum_n h^n_C \), that
\[
E(s, t) = \sum_{n \geq 0} [k^n_C h(\cdot, t)](s) = h(s, t) + \int_0^s \kappa_C(s, v)h(v, t)dv,
\]
with a uniformly bounded kernel \( \kappa_C \). Plugging \( \text{(4.1)} \) into \( \text{(1.18)} \), we find by Fubini’s theorem that
\[
\chi(s, t) = s \wedge t + \int_0^s \int_0^{t \wedge u} \nu'(C(v, u))du\kappa_1(s, v)dv + \int_0^s \chi(v, t)\kappa_2(s, v)dv,
\]
for some uniformly bounded functions \( \kappa_1 \) and \( \kappa_2 \) which depend only on \( C \) and \( D \). Applying Picard’s iterations once more, now with respect to the integral operator \( [\kappa_2 g](s) = \int_0^s \kappa_2(s, v)g(v)dv \), we deduce that for some uniformly bounded \( \kappa_3 \) and \( \kappa_4 \),
\[
\chi(s, t) = s \wedge t + \int_0^s [(u \wedge t)\kappa_3(s, u) + \int_0^{t \wedge u} \nu'(C(u, v))dv] \kappa_4(s, u)du.
\]
On \( s \geq t \), the function \( s \wedge t = t \) is continuously differentiable, hence we easily conclude by Fubini’s theorem that \( t \to \chi(s, t) \) is continuously differentiable on \( s \geq t \), with \( \chi(s, t) = \int_0^t R(s, u)du \) for the bounded continuous function
\[
R(s, t) = 1 + \int_t^s [\kappa_3(s, u) + \nu'(C(u, t))] \kappa_4(s, u)du.
\]
In particular, \( R(s, s) = 1 \) for all \( s \). The condition \( E(s, t) = E(s, s) \) for \( s \geq t \) implies by \( \text{(1.18)} \) that \( \chi(s, t) = \chi(s, s) \) for \( s \geq t \). Similarly, with \( E(s, 0) = 0 \), it follows that \( \chi(s, 0) = 0 \) for all \( s \). From \( \text{(1.17)} \), we have that \( C(s, t) - \chi(s, t) \) is differentiable with respect to its second argument \( t \), with a bounded, continuous derivative \( D = \partial_2(C - \chi) \). Consequently, \( \partial_2 C = D + R \) where \( R(s, t) = (\partial_2 \chi)(s, t) = 0 \) for all \( t > s \) due to the boundary condition \( \chi(s, t) = \chi(s, s) \). Further, \( C(s, t) = C(t, s) \) implying that \( \partial_1 C(s, t) = \partial_2 C(t, s) = D(t, s) + R(t, s) \) on \([0, T]^2\). Thus, combining the identity
\[
C(s, t \vee s)\nu'(C(t \vee s, t)) - C(s, 0)\nu'(C(0, t)) = \int_0^{t \wedge s} \nu'(C(u, u))(\partial_2 C)(s, u)du
\]
\[
+ \int_0^{t \wedge s} C(s, u)\nu''(C(t, u))(\partial_2 C)(t, u)du,
\]
with \( \text{(1.19)} \) we have that for all \( t, s \in [0, T]^2 \),
\[
D(s, t) = -f'(K(t))C(t, s) + \int_0^{t \wedge s} \nu'(C(t, u))R(s, u)du + \int_0^{t \wedge s} C(s, u)\nu''(C(t, u))R(t, u)du.
\]
Interchanging \( t \) and \( s \) in \( \text{(4.2)} \) and adding \( R(t, s) = 0 \) when \( s > t \), results for \( s > t \) with
\[
(\partial_1 C)(s, t) = -f'(K(s))C(s, t) + \int_0^s \nu'(C(s, u))R(t, u)du + \int_0^s C(t, u)\nu''(C(s, u))R(s, u)du,
\]
which is \( \text{(1.13)} \) for \( \beta = 1 \).
Since \( K(s) = C(s, s) \), with \( C(s, t) = C(t, s) \) and \( \partial_2 C = D + R \), it follows that for all \( h > 0 \),
\[
K(s) - K(s - h) = \int_{s-h}^s (D(s, u) + R(s, u))du + \int_{s-h}^s (D(s - h, u) + R(s - h, u))du.
\]
Recall that \( R(s,u) = 0 \) for \( u > s \), hence, dividing by \( h \) and taking \( h \downarrow 0 \), we thus get by the continuity of \( D \) and that of \( R \) for \( s \geq t \) that \( K(\cdot) \) is differentiable, with \( \partial_s K(s) = 2D(s,s)+R(s,s) = 2D(s,s)+1 \), resulting by (1.12) with (1.14) for \( \beta = 1 \).

Further, it follows from (1.18) that \( (\partial_1 \chi)(u,t) = E(u,t) + 1_{u<t} \). Hence, combining the identity

\[
\chi(s,t)\nu'(C(s,s)) - \chi(0,t)\nu'(C(s,0)) = \int_0^s \nu'(C(s,u))(\partial_1 \chi)(u,t)du + \int_0^s \chi(u,t)\nu''(C(s,u))(\partial_2 C)(s,u)du,
\]

with (1.20) we have that for all \( T \geq s \geq t \),

\[
E(s,t) = -f'(K(s))\chi(s,t) + \int_0^s \chi(u,t)\nu''(C(s,u))R(s,u)du
\]
(recall that \( \chi(0,t) = 0 \)). Let

\[
g(s,t) := -f'(K(s))R(s,t) + \int_0^s R(u,t)\nu''(C(s,u))R(s,u)du,
\]
for \( s,t \in [0,T]^2 \). Recall that \( \chi(s,t) = \int_0^t R(s,v)dv \), so by Fubini’s theorem, (4.3) amounts to

\[
E(s,t) = \int_0^t g(s,v)dv
\]
for all \( s,t \leq T \). Putting this into (1.18) we have by yet another application of Fubini’s theorem that

\[
\int_0^t R(s,u)du = \chi(s,t) = t + \int_0^s \int_0^{t\wedge u} g(u,v)dvdu = t + \int_0^t \int_v^s g(u,v)dvdu,
\]
for any \( s \geq t \). Consequently, for every \( t \leq s \),

\[
R(s,t) = 1 + \int_t^s g(u,t)du,
\]

implying that \( \partial_1 R = g \) for a.e. \( s > t \), which in view of (1.5) gives (1.12) for \( \beta = 1 \), thus completing the proof of the lemma.

We proceed by showing that the system of equations of interest to us admits at most one solution.

**Proposition 4.2.** Let \( T \geq 0 \) and \( \Delta_T = \{ s,t \in (\mathbb{R}^+)^2 : 0 \leq t \leq s \leq T \} \). There exists at most one solution \((R,C,K) \in C^1_b(\Delta_T) \times C^1_b(\Delta_T) \times C^1_b([0,T])\) to (S):= (1.1, (1.3), (1.4)) with \( C(s,t) = C(t,s) \) and boundary conditions

\[
R(s,s) = 1 \quad \forall s \geq 0
\]

\[
C(s,s) = K(s) \quad \forall s \geq 0
\]

\[
C(0,0) = K(0) \quad \text{known.}
\]

**Proof:** As mentioned already, we may and shall take \( \beta = 1 \) (by scaling \( \beta \nu \mapsto \nu \), with \( \beta^2 \psi \mapsto \psi \) accordingly), just as we have done throughout this paper.

Assume that there are two solutions \((R,C,K)\) and \((\tilde{R},\tilde{C},\tilde{K})\) of (S) with boundary condition \((BC):= (1.5)\). We shall prove by Gronwall’s type argument that these two solutions have to coincide. To do so we first show that the response function \( R \) is a Lipschitz function of the covariance functions \( (C,K) \) and then that the covariances obey integro-differential Gronwall type inequalities. We then use Gronwall arguments repeatedly to conclude. In what follows, \( T \) is fixed and all the constants (which eventually depend on \( T \)) will be denoted by \( M \), even though they may change from line to line.
• $R$ is a Lipschitz function of the covariance

Let $(R, C, K)$ be a solution to (S) and denote

$$H_C(s, t) := e^{\int_t^s f'(K(u))du}R(s, t).$$

Then, from (1.12) and (1.15), we deduce that $H$ satisfies

$$\partial_s H_C(s, t) = \int_t^s H_C(u, t)H_C(s, u)\nu''(C(s, u))du, \text{ for } s \geq t, \quad H(t, t) = 1.$$ 

Note that $H$ only depends on $C$. This equation was studied in [16] where it was shown that if $NC_n$ denotes the set of involutions of $\{1, \cdots, 2n\}$ without fixed points and without crossings and if $\text{cr}(\sigma)$ is defined to be the set of indices $1 \leq i \leq 2n$ such that $i < \sigma(i)$, $H$ is given by

$$H_C(s, t) = 1 + \sum_{n \geq 1} \sum_{1 \leq i_1 \leq \cdots \leq i_n \leq s} \prod_{\sigma \in NC_n} \nu''(C(t_{\sigma(i)}, i))dt_1 \cdots dt_n$$

The number of non-crossing partitions $|NC_n|$ of $2n$ elements is given by the Catalan number $C_n$ which is at most $2^n$. As $\int_{t_1 \leq i_1 \leq \cdots \leq t_n \leq t} dt_1 \cdots dt_n \leq M/(2n)!$ and $\sup_{t, s} \nu''(C(s, t))$ is uniformly bounded by hypothesis, the above sum converges absolutely. Further, telescoping $\prod_{\sigma \in \text{cr}(\sigma)} \nu''(C(t_{\sigma(i)}, i)) = \prod_{\sigma \in \text{cr}(\sigma)} \nu''(C(t_{\sigma(i)}, i))$, by the uniform Lipschitz continuity of $\nu''(\cdot)$ on compacts, we thus find a finite constant $M$ such that for any pair $C, \tilde{C} \in \mathcal{C}_b(\Delta_T)$ and any $t, s \in \Delta_T$, $\tilde{C} \in \mathcal{C}_b(\Delta_T)$ and any $t, s \in \Delta_T$,

$$|H_C(s, t) - H_{\tilde{C}}(s, t)| \leq M \int_{t \leq t_2 \leq t_1 \leq s} |C(t_1, t_2) - \tilde{C}(t_1, t_2)|dt_1 dt_2.$$

Thus, if $(R, C, K)$ and $(\tilde{R}, \tilde{C}, \tilde{K})$ are two solutions of (S) in $\mathcal{C}_b^1(\Delta_T) \times \mathcal{C}_b^1(\Delta_T) \times \mathcal{C}_b^1([0, T])$, since $K$ is uniformly bounded and $f'(\cdot)$ is locally Lipschitz, we obtain

$$|R(s, t) - \tilde{R}(s, t)| \leq M \int_{t_2 \leq t_1 \leq t} |C(t_1, t_2) - \tilde{C}(t_1, t_2)|dt_1 dt_2 + M \int_t^s |K(u) - \tilde{K}(u)|du.$$

• Bounds on the difference of the covariances on $s \geq t$

Integrating (1.13) yields for $s \geq t$

$$C(s, t) = K(t) - \int_t^s f'(K(u))C(u, t)du + \int_t^s du \int_t^s \nu''(C(u, v))R(t, v)$$

$$+ \int_t^s du \int_t^u \nu''(C(u, v))C(t, v)R(u, v) + \int_t^s du \int_t^u \nu''(C(u, v))C(v, t)R(u, v).$$

Hence, if $(R, C, K)$ and $(\tilde{R}, \tilde{C}, \tilde{K})$ are two solutions of (S),

$$|C - \tilde{C}|(s, t) \leq M \left[ |K - \tilde{K}|(t) + \int_t^s |K - \tilde{K}|(u)du + \int_t^s |C - \tilde{C}|(u, t)du + \int_t^s du \int_0^s dv|C - \tilde{C}|(u, v)$$

$$+ \int_t^s du \int_0^s dv|C - \tilde{C}|(t, v) + \int_t^s du \int_0^s dv|R - \tilde{R}|(t, v) + \int_t^s du \int_0^s dv|R - \tilde{R}|(u, v)$$

$$+ \int_t^s du \int_0^u dv|C - \tilde{C}|(u, v) + \int_t^s du \int_t^u dv|C - \tilde{C}|(v, t) + \int_t^s du \int_t^u dv|R - \tilde{R}|(u, v) \right]$$

$$= I_1(s, t) + I_2(s, t) + \cdots + I_{10}(s, t)$$

(4.10)
• **Bounds on the differences of the covariances on the diagonal**

Similarly, integrating (1.14) gives

\[
K(t) = K(0) - 2 \int_0^t f'(K(u))K(u)du + t + 2 \int_0^t du \int_0^u dv \psi(C(u, v))R(u, v),
\]

yielding in case \(K(0) = \tilde{K}(0)\) that

\[
(4.1) \quad |K - \tilde{K}|(t) \leq M \left[ \int_0^t |K - \tilde{K}|(u)du + \int_0^t du \int_0^u dv |C - \tilde{C}|(u, v)dv + \int_0^t du \int_0^u dv |R - \tilde{R}|(u, v)dv \right]
\]

Plugging (4.9) into (4.11) yields

\[
(4.12) \quad |K - \tilde{K}|(t) \leq M \left[ \int_{0 \leq t_1 \leq t_2 \leq t} |C - \tilde{C}|(t_2, t_1)dt_1dt_2 + \int_0^t |K - \tilde{K}|(u)du \right]
\]

Recall that by Gronwall’s lemma, if \(h, g\) are two non-negative functions such that

\[
h(t) \leq g(t) + A \int_0^t h(s)ds
\]

for some \(A \geq 0\), then

\[
h(t) \leq g(t) + A \int_0^t g(v)e^{A(t-v)}dv \leq e^{At}g(t)
\]

where the last inequality holds when \(g\) is non-decreasing. Applying this inequality with

\[
g(t) = \int_{0 \leq t_1 \leq t_2 \leq t} |C - \tilde{C}|(t_2, t_1)dt_1dt_2
\]

which is non-negative and non-decreasing yields

\[
(4.13) \quad |K - \tilde{K}|(t) \leq M \left[ \int_{0 \leq t_1 \leq t_2 \leq t} |C - \tilde{C}|(t_2, t_1)dt_1dt_2 \right]
\]

• **The final fixed point argument**

We now consider

\[
D(s) := \int_0^s |C - \tilde{C}|(s, t)dt,
\]

noting that (4.9) and (4.13) imply that

\[
(4.14) \quad |R(s, t) - \tilde{R}(s, t)| \leq M \int_{0 \leq t_2 \leq t_1 \leq t} |C(t_1, t_2) - \tilde{C}(t_1, t_2)|dt_1dt_2 = M \int_{0 \leq t_1 \leq s} D(t_1)dt_1,
\]

and

\[
(4.15) \quad |K - \tilde{K}|(t) \leq M \left[ \int_{0 \leq t_2 \leq t} D(t_2)dt_2 \right].
\]

Thus, integrating (4.10) with respect to \(t\) and observing that

\[
\int_0^s (I_1(s, t) + I_2(s, t))dt \leq M \int_0^s D(u)du \quad \text{by (4.15)},
\]

\[
\int_0^s (I_4(s, t) + I_8(s, t))dt \leq M \int_0^s dt \int_t^s du \int_0^u dv |C - \tilde{C}|(u, v) \leq M \int_0^s D(u)du \quad \text{and}
\]

\[
\int_0^s (I_9(s, t) + I_5(s, t) + I_3(s, t))dt \leq M \int_0^s D(u)du \quad \text{by definition of} \ D \text{and Fubini},
\]

\[
\int_0^s (I_6(s, t) + I_7(s, t) + I_{10}(s, t))dt \leq M \int_0^s dt \int_t^s du \int_0^u dv |R - \tilde{R}|(u, v) \leq M \int_0^s D(u)du \quad \text{by (4.14)},
\]

\[
\int_0^s (I_0(s, t) + I_{11}(s, t))dt \leq M \int_0^s dt \int_t^s du \int_0^u dv |D(u) - \tilde{D}(u)| \leq M \int_0^s D(u)du.
\]
we obtain from (4.10) that

\[ D(s) \leq M \int_0^s D(u)du. \]

Recall that \( D \) is non-negative and non-decreasing, so by the preceding Gronwall argument, now with \( g = 0 \) we conclude that \( D(s) = 0 \) for all \( s \in [0,T] \). This in turn implies by (4.14) and (4.15) that

\[ K(t) = \tilde{K}(t), \quad R(s,t) = \tilde{R}(s,t) \quad \text{for all } (t,s) \in \Delta_T \]

and \( C(s,t) = \tilde{C}(s,t) \) for almost all \( t \leq s \) and all \( s \leq T \). Either by (4.10) or directly by the continuity of the covariances we conclude that

\[ C(s,t) = \tilde{C}(s,t) \quad \text{for all } (t,s) \in \Delta_T \]

which finishes the proof. \( \square \)

We conclude this section with the,

**Proof of Theorem 1.2**: Recall Proposition 2.3 that we have pre-compactness of \((A^a_N, F^a_N, \chi^a_N, C^a_N) : [0,T]^2 \rightarrow \mathbb{R}^4\), in the topology of uniform convergence on \([0,T]^2\). This implies the existence of limit points for this sequence. By Proposition 1.3 every such limit point is a solution of the integral equations (1.17)–(1.20) with the stated symmetry and boundary conditions. By Lemma 4.1 each such solution results with \( C \) and \( \chi \) (i.e. \( R \)) that satisfy the integro-differential equations (1.12)–(1.14). In view of Proposition 4.2 the latter system admits at most one solution per given boundary conditions. Hence, we conclude that the sequence \((\chi^a_N, C^a_N)\) converges uniformly in \([0,T]^2\) to the unique solution of (1.12)–(1.14) subject to the appropriate boundary conditions. Further, by (2.19) of Proposition 2.4 both \( C_N - C^a_N \) and \( \chi_N - \chi^a_N \) converge uniformly to zero, almost surely. Thus, the solution of (1.12)–(1.14) is also the unique almost sure uniform (in \( s,t \)) limit of \((\chi_N, C_N)\), as stated in Theorem 1.2. The \( L_p \) convergence then follows by the uniform bounds on moments of \( C_N \) and \( \chi_N \) (see Proposition 2.3), thus completing the proof of the theorem. \( \square \)

**Appendix A. Itô’s calculus**

Let \( \{x^i_t, y^i_t, z^i_t, w^i_t\}_{t \geq 0, i \in \mathbb{N}} \) be semi-martingales such that,

\[ d(r^p_t, p^i_t) = \delta_{i-j}dt \]

for any \( p, r \in \{x, y, z, w\} \). Denoting, for \( p, r \in \{x, y, z, w\}, s, t \geq 0, N \in \mathbb{N}, \)

\[ K^N_{p,r}(s,t) := \frac{1}{N} \sum_{i=1}^N p^i_r z^i_t, \]

we already made use of the following simple stochastic calculus lemma.
Lemma A.1. For any polynomials $P, Q$, and any $\sigma, \theta, v \geq 0$,

\[
P(K_{x,y}^N(\sigma, \theta))Q(K_{z,w}^N(\sigma, v)) = P(K_{x,y}^N(0, \theta))Q(K_{z,w}^N(0, v))
\]

\[
+ \int_0^\sigma P'(K_{x,y}^N(u, \theta))Q(K_{z,w}^N(u, v))du
\]

\[
+ \int_0^\sigma P(K_{x,y}^N(u, \theta))Q'(K_{z,w}^N(u, v))du
\]

\[
+ \frac{1}{2N}K_{y,y}^N(\theta, \theta)\int_0^\sigma P''(K_{x,y}^N(u, \theta))Q(K_{z,w}^N(u, v))du
\]

\[
+ \frac{1}{2N}K_{w,w}^N(v, v)\int_0^\sigma P(K_{x,y}^N(u, \theta))Q''(K_{z,w}^N(u, v))du
\]

\[
+ \frac{1}{N}K_{y,w}^N(\theta, v)\int_0^\sigma P'(K_{x,y}^N(u, \theta))Q'(K_{z,w}^N(u, v))du
\]

where

\[
d_uK_{z,w}^N(u, v) := \frac{1}{N} \sum_{i=1}^N w^i d z^i_u,
\]

and all the stochastic integrals are defined via our convention (of putting terms such as $y^i_{\theta} \cdots y^i_{\theta} w^j_{v} \cdots w^j_{v}$ outside the integral).

Proof: By the bi-linearity of the formula given, it is enough to prove the lemma for $P(x) = x^a$ and $Q(x) = x^b$. In this case, writing

\[
(K_{x,y}^N(\sigma, \theta))^a (K_{z,w}^N(\sigma, v))^b = N^{-(a+b)} \sum_{i_1,\ldots,i_a,j_1,\ldots,j_b} y^i_{\theta} \cdots y^i_{\theta} w^j_{v} \cdots w^j_{v} \left( x^{i_1}_{\sigma} \cdots x^{i_a}_{\sigma} z^{j_1}_{\sigma} \cdots z^{j_b}_{\sigma} \right),
\]

and using Itô’s formula for $x^{i_1}_{\sigma} \cdots x^{i_a}_{\sigma} z^{j_1}_{\sigma} \cdots z^{j_b}_{\sigma}$ gives the stated result. \(\square\)

Appendix B. Supremum of Gaussian processes indexed on large dimensional spheres.

In this section we prove the bound \([2,12]\) which is a direct consequence of the following general lemma about supremum of Gaussian processes indexed on large dimensional balls. The outline for a direct proof of such a result by a chaining argument was kindly communicated to us by A. Bovier, whom we thank gratefully. This chaining argument can be adapted rather straightforwardly from \([2]\). Anton Bovier also mentioned that this result should be the consequence of a more general one. We have indeed found the proper way to see it as a consequence of classical and well known facts on Gaussian processes, and to give simple references.

Let $(X_N(x))$ be a sequence of real valued, centered Gaussian processes indexed by $x \in \mathbb{R}^N$. Consider, for every $\rho > 0$, the closed Euclidean ball $B_N(0, \rho)$ in $\mathbb{R}^N$, and define

\[
X_N^*(\rho) = \sup_{x \in B_N(0, \rho)} \frac{|X_N(x)|}{\sqrt{N}}
\]

We will also introduce the usual metric on $\mathbb{R}^N$ associated to the process $X_N$,

\[
d_X(x, y) = \mathbb{E}[|X_N(x) - X_N(y)|^2]^{1/2}
\]

We denote by $\| \cdot \|$ the Euclidean norm and by $(x, y)_N$ the corresponding inner product on $\mathbb{R}^N$. 

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Lemma B.1. Suppose that
\[ \sup_N \mathbb{E}[X_N(0)^2] < \infty \] (B.1)
and that
\[ \sup_N \sup_{x,y \in B_N(0,\rho)} \frac{d_X(x,y)}{\|x - y\|} < \infty. \] (B.2)
Then, for every \( k \in \mathbb{N} \)
\[ \sup_N \mathbb{E}[X_N(\rho)^k] < \infty. \] (B.3)
Moreover, there exists a constant \( \kappa < \infty \) such that for all \( N \) and every \( t > 0 \),
\[ \mathbb{P}[X_N(\rho) \geq \kappa + t] \leq \exp(-Nt^2/\kappa). \] (B.4)

Proof: This result is a direct consequence of Dudley's theorem (13). Indeed, the assumption (B.2) implies that for any \( N \) and \( \epsilon > 0 \) one can cover \( B_N(0,\rho) \) by the union of certain \( C(\rho) e^{-N} \) balls of radius \( \epsilon \), in the metric \( d_X \), where the constant \( C(\rho) \) depends on \( \rho \) but not on the dimension \( N \). Thus, Dudley’s theorem (see also [13, Theorem 11.17]) shows that
\[ \mathbb{E}\left[ \sup_{x \in B_N(0,\rho)} X_N(x) \right] \leq C'(\rho) \sqrt{N}, \]
where the constant \( C'(\rho) \) again depends only on \( \rho \) and not on \( N \). Using the obvious fact that
\[ \mathbb{E}\left[ \sup_{x \in B_N(0,\rho)} |X_N(x)| - \mathbb{E}|X_N(0)| \right] \leq \mathbb{E}\left[ \sup_{x,y \in B_N(0,\rho)} |X_N(x) - X_N(y)| \right] \]
\[ = \mathbb{E}\left[ \sup_{x,y \in B_N(0,\rho)} \{X_N(x) - X_N(y)\} \right] \leq 2 \mathbb{E}\left[ \sup_{x \in B_N(0,\rho)} X_N(x) \right], \]
and the assumption (B.1), we see that the conclusion (B.3) holds for \( k = 1 \). Thus, \( X_N \) admits a version with almost all sample paths bounded and uniformly continuous on \( B_N(0,\rho) \). One can then consider \( X_N \) as an (infinite dimensional) Gaussian vector in the space of continuous functions on the ball \( B_N(0,\rho) \), equipped with the supremum norm. It is also a well known fact that, for such a Gaussian vector, all moments of the norm are controlled by the first (e.g. see the last statement of [13, Corollary 3.2]). This is thus enough to ensure that (B.3) holds for every \( k \in \mathbb{N} \).

The tail estimate (B.4) is also classical in the Gaussian context. For instance, the assumptions (B.1) and (B.2) immediately imply that the weak variance
\[ \sigma(X_N) = \sup_{x \in B_N(0,\rho)} \left\{ \mathbb{E}[X_N(x)^2]^{1/2} \right\} \]
of [13, page 56] is bounded in \( N \). Hence, by [13] estimate (3.2), page 57] it is easy to see that there exists a finite constant \( \kappa > \sup_N \mathbb{E}|X_N(\rho)| \) for which (B.4) applies.

We proceed to apply Lemma B.1 to the situation of interest here. To this end, fixing \( p \in \mathbb{N} \) consider the Gaussian process defined on \( (\mathbb{R}^N)^p \)
\[ X_{N,\rho}(x) = \sum_{1 \leq i_j \leq N, 1 \leq j \leq p} G_{\{i_1, \ldots, i_p\}} x_{i_1}^1 x_{i_2}^2 x_{i_3}^3 \cdots x_{i_p}^p \]
with independent, centered Gaussian variables \( G_{\{i_1, \ldots, i_p\}} \) of variances \( c(\{i_1, \ldots, i_p\}) \) of (13). Considering in \( (\mathbb{R}^N)^p \) the Cartesian product
\[ B(\rho) = \prod_{i=1}^p \left\{ x^i \in \mathbb{R}^N, \|x^i\| \leq \rho \right\} \]
of $p$ Euclidean balls, we wish to estimate the moments and tail of $\sup_{x \in B(\rho)} |X_{N,p}(x)|$. To this end, note that since $X_{N,p}$ is a symmetric $p$-linear form on $\mathbb{R}^N$ (i.e., $X_{N,p}(x)$ is invariant to permutations of the vectors $x^1, \ldots, x^p$), we have by polarization that

$$\sup_{x \in B(\rho)} |X_{N,p}(x)| \leq C(p) \sup_{\|u\| \leq \rho p} |X_N(u)|,$$

where one defines $X_N(u)$ for $u \in \mathbb{R}^N$ by

$$X_N(u) = X_{N,p}(u, \ldots, u) = \sum_{1 \leq i_j \leq N} G_{(i_1, \ldots, i_p)} u_{i_1}u_{i_2} \cdots u_{i_p}.$$

The Gaussian process $X_N(u)$ obviously satisfies (B.1) since $X_N(0) = 0$ for all $N$. Turning to check that (B.2) is satisfied as well, note that

$$X_N(u) = N^{(p-1)/2} \sum_{i=1}^N G^i(u)u_i,$$

for the case where $\nu'(r) = (p-1)!r^{p-1}$. Thus, by (3.3) and (3.7), the covariance of the process $X_N$ is

$$\mathbb{E}[X_N(u)X_N(v)] = p!(u,v)_N^p,$$

so that

$$d_X(u,v)^2 = p!\left[ (u,v)_N^p + (v,v)_N^p - 2(u,v)_N^p \right].$$

It is then easy to check that for all $u$ and $v$,

$$d_X(u,v) \leq C(p)\|u - v\|,$$

yielding that (B.2) holds. Hence, by Lemma B.1 for every positive integers $p$ and $k$,

$$(B.6) \quad \sup_N \mathbb{E}\left[ \sup_{x \in B(\rho)} \left( \frac{|X_{N,p}(x)|}{\sqrt{N}} \right)^k \right] \leq \sup_N \mathbb{E}[X_N^k] < \infty.$$

In view of (2.1), combining (B.6) for $p = 1, \ldots, m$ then proves the bound (2.12). We further have by (B.4) and (B.5) that for some $\tilde{\kappa} < \infty$, any $p \leq m$, all $N$ and every $t > 0$,

$$\mathbb{P}\left[ \sup_{x \in B(\rho)} \left( \frac{|X_{N,p}(x)|}{\sqrt{N}} \right) \geq \tilde{\kappa} + t \right] \leq \exp(-Nt^2/2\tilde{\kappa}),$$

from which it follows that,

$$(B.7) \quad \mathbb{P}[\|J\|_\infty \geq \tilde{\kappa} + t] \leq m \exp(-Nt^2/2\tilde{\kappa}).$$

Appendix C. Gaussian change of measure identities

The change of measure that is key to the proof of Lemma 3.3 is a special case of the following “standard” Gaussian computation (compare (E.16) and (C.1)).

**Proposition C.1.** Suppose under the law $\mathbb{P}$ we have a finite collection $J = \{J_\alpha\}_\alpha$ of non-degenerate, independent, centered Gaussian random variables, and $G_s^i = \sum_\alpha J_\alpha L_s^i(\alpha)$ for all $s \in [0, \tau]$ and $i \leq N$, where for each $\alpha$ the coefficients $L_s^i$ which are independent of $J$ are in $L_2(\{1, \ldots, N\} \times [0, \tau])$. Suppose further that $U_s^i$ is a continuous semi-martingale, independent of $J$ and such that for each $\alpha$ the stochastic integral

$$\mu_\alpha := \sum_{i=1}^N \int_0^\tau L_s^i(\alpha)dU_s^i,$$
is well defined and almost surely finite. Let $\mathbb{P}^*$ denote the law of $J$ such that $\mathbb{P}^* = L_\tau/\mathbb{E}(L_\tau)\mathbb{P}$, where

\[
L_\tau = \exp \left\{ \sum_{i=1}^{N} \int_{0}^{\tau} G_i^s dU^i_s - \frac{1}{2} \sum_{i=1}^{N} \int_{0}^{\tau} (G_i^s)^2 ds \right\}.
\]

Let $k_{ts}^{ij} = \mathbb{E}(G_t^s G_s^i)$, $V_s^i = \mathbb{E}^*(G_s^i)$ and $\Gamma_{ts}^{ij} = \mathbb{E}^*[(G_t^s - V_t^i)(G_s^i - V_s^j)]$. Then, for any $s \leq \tau$ and $i \leq N$,

\[
V_s^i + [k_\tau V]^i_s = [k_\tau \circ dU]_s^i,
\]

and for any $s, t \leq \tau$ and $i, l \leq N$,

\[
\sum_{j=1}^{N} \int_{0}^{\tau} k_{su}^{ij} \Gamma_{ut}^{jl} du + \Gamma_{st}^{il} = k_{st}^{il}.
\]

**Proof:** Let $v_\alpha = \mathbb{E}(J_\alpha^2) > 0$ denote the variance of $J_\alpha$ and

\[
R_{\alpha\gamma} := \sum_{i=1}^{N} \int_{0}^{\tau} L_i^\alpha(\alpha) L_i^\gamma(\gamma) du,
\]

observing that

\[ L_\tau = \exp \left\{ \sum_\alpha J_\alpha \mu_\alpha - \frac{1}{2} \sum_{\alpha, \gamma} J_\alpha J_\gamma R_{\alpha\gamma} \right\}. \]

With $D = \text{diag}(v_\alpha)$ a positive definite matrix and $R = \{R_{\alpha\gamma}\}$ positive semi-definite, it follows from this representation of $L_\tau$ that under $\mathbb{P}^*$ the random vector $J$ has a Gaussian law with covariance matrix $(D^{-1} + R)^{-1}$ and mean vector $q = \{q_\alpha\} = (D^{-1} + R)^{-1} \mu$. Hence, for any $\alpha$,

\[
q_\alpha + v_\alpha \sum_\gamma R_{\alpha\gamma} q_\gamma = v_\alpha \mu_\alpha.
\]

As $k_{su}^{ij} = \sum_\alpha L_i^\alpha(\alpha) v_\alpha L_s^j(\alpha)$, it is not hard to check that

\[
[k_\tau \circ dU]_s^i := \sum_{j=1}^{N} \int_{0}^{\tau} k_{su}^{ij} dU^j_s = \sum_\alpha L_i^\alpha(\alpha) v_\alpha \mu_\alpha.
\]

Obviously, $V_s^i = \sum_\alpha L_s^i(\alpha) q_\alpha$ so we get (C.2) out of (C.5) upon verifying that

\[
[k_\tau V]^i_s := \sum_{j=1}^{N} \int_{0}^{\tau} k_{su}^{ij} V^j_s du = \sum_\alpha L_i^\alpha(\alpha) v_\alpha q_\gamma \sum_{j=1}^{N} \int_{0}^{\tau} L_s^j(\gamma) du = \sum_\alpha L_i^\alpha(\alpha) v_\alpha R_{\alpha\gamma} q_\gamma,
\]

with the last identity due to (C.4).

Turning to prove (C.3), since $\Gamma_{ut}^{jl}$ is the covariance of $G_t^i$ and $G_t^j$ under the tilted law $\mathbb{P}^*$, we have that

\[
\Gamma_{ut}^{jl} = \sum_{\alpha, \gamma} L_i^\alpha(\alpha) [(D^{-1} + R)^{-1}]_{\alpha\gamma} L_j^\gamma(\gamma),
\]

and hence by (C.4) we see that

\[
\sum_{j=1}^{N} \int_{0}^{\tau} k_{su}^{ij} \Gamma_{ut}^{jl} du = \sum_{\alpha, \gamma} L_i^\alpha(\alpha) v_\alpha [R(D^{-1} + R)^{-1}]_{\alpha\gamma} L_j^\gamma(\gamma).
\]

With $D = \text{diag}(v_\alpha)$ we easily get (C.3) out of the matrix identity $(I + DR)(D^{-1} + R)^{-1} = D$. \qed
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COURANT INSTITUTE, 251 MERCER STREET, NEW YORK, NY 10012, AND MATHEMATICS, EPFL, 1015 LAUSANNE, SWITZERLAND

E-mail address: beunarous@cims.nyu.edu

DEPARTMENT OF STATISTICS AND DEPARTMENT OF MATHEMATICS, STANFORD UNIVERSITY, STANFORD, CA 94305.

E-mail address: amir@math.stanford.edu

UMPA, ECOLE NORMALE SUPERIEURE DE LYON, 46 ALLÉE D’ITALIE, 69364 LYON CEDEX 07, FRANCE

E-mail address: Alice.Guionnet@umpa.ens-lyon.fr