NEW INTEGRAL REPRESENTATIONS FOR THE FOX-WRIGHT FUNCTIONS
AND ITS APPLICATIONS II

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Abstract. In this paper our aim is to establish new integral representations for the Fox–Wright function
\( pFq \) when

\[
\mu = \sum_{j=1}^{q} \beta_j - \sum_{k=1}^{p} \alpha_k + \frac{p-q}{2} = -m, \quad m \in \mathbb{N}_0.
\]

In particular, closed-form integral expressions are derived for the four parameter Wright function under
a special restriction on parameters. Exponential bounding inequalities are derived for a class of the
Fox–Wright function. Moreover, complete monotonicity property is presented for these functions.

1. Introduction

We use a definition of the Fox–Wright (generalized hypergeometric) function by its series

\[
\Psi_p Q_{\beta_1, \ldots, \beta_p} (\alpha_1, \ldots, \alpha_p; z) = \sum_{k=0}^{\infty} \frac{\prod_{i=1}^{p} \Gamma(\alpha_i + kA_i)}{\prod_{j=1}^{q} \Gamma(\beta_j + kB_j) \cdot k!} z^k,
\]

where, as usual,

\[
\mathbb{N} = \{1, 2, 3, \ldots\}, \quad \mathbb{N}_0 = \mathbb{N} \cup \{0\},
\]

\( \mathbb{R}, \mathbb{R}_+ \) and \( \mathbb{C} \) stand for the sets of real, positive real and complex numbers, respectively. This function
was first introduced by Wright [7] in 1935, who also derived some of its important properties including
asymptotic behavior.

The convergence conditions and convergence radius of the series at the right-hand side of (1.1) immediately follow from the known asymptotic of the Euler Gamma-function. To formulate the results, let us first introduce the following notations:

\[
\Delta = \sum_{j=1}^{q} B_j - \sum_{i=1}^{p} A_i, \quad \rho = \left( \prod_{i=1}^{p} A_i^{-\alpha_i} \right) \left( \prod_{j=1}^{q} B_j^{B_j} \right), \quad \mu = \sum_{j=1}^{q} \beta_j - \sum_{k=1}^{p} \alpha_k + \frac{p-q}{2}
\]

The defining series in (1.1) converges in the whole complex \( z \)-plane if \( \Delta > -1 \). If \( \Delta = -1 \), then the series
in (1.1) converges for \( |z| < \rho \), and \( |z| = \rho \) under the condition \( \Re(\mu) > \frac{1}{2} \), see [15] for details. If, in the
definition (1.1), we set

\[
A_1 = \ldots = A_p = 1 \quad \text{and} \quad B_1 = \ldots = B_q = 1,
\]

we get the relatively more familiar generalized hypergeometric function \( pFq \) given by

\[
pFq_{\alpha_1, \ldots, \alpha_p} (\beta_1, \ldots, \beta_q; z) = \sum_{k=0}^{\infty} \frac{\prod_{i=1}^{p} \Gamma(\alpha_i)}{\prod_{j=1}^{q} \Gamma(\beta_j + kB_j) \cdot k!} (z)^k, \quad (\alpha_j > 0, \beta_j \notin \mathbb{Z}_0^+).
\]

Moreover, both the Wright function \( W_{\alpha, \beta}(z) \) and the Mittag-Leffler function \( E_{\alpha, \beta}(z) \), are particular cases of the Fox–Wright function (1.1):

\[W_{\alpha, \beta}(z) = e^{\Psi_{1} (z, \alpha)} = 1, \quad E_{\alpha, \beta}(z) = 1 \Psi_{1} (1, \alpha) \]

Note important properties for this functions including its Turán, Lazarević and Wilker type inequalities,
was proved by Mehrez [11] and Mehrez et al in [9–10].

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\[ \text{NEW INTEGRAL REPRESENTATIONS FOR THE FOX-WRIGHT FUNCTIONS} \]
In a recent papers [12, 13, 14], the author have studied certain advanced properties of the Fox-Wright function including its new integral representations, the Laplace and Stieltjes transforms, Luke inequalities, Turán type inequalities and completely monotonicity property are derived. In particular, it was shown there that the following Fox-Wright functions are completely monotone:

\[ pΨq \left[ \frac{z}{(β_q, A_q)} \right] = \psi^0, \quad z > 0, \]

\[ p+1Ψq \left[ \frac{z}{(β_q, A_q)} \right] = \psi^0, \quad z > 0, \]

and has proved that the Fox’s H-function \( H_{q,p}^{p,0} \) constitutes the representing measure for the Fox-Wright function \( pΨq \), if \( p > 0 \), i.e., [12, Theorem 1]

(1.4)

\[ pΨq \left[ \frac{z}{(β_q, A_q)} \right] = \int_0^1 e^{zt} H_{q,p}^{p,0} \left( t \frac{z}{(β_q, A_q)} \right) \frac{dt}{t}. \]

when \( μ > 0 \). Here, and in what follows, we use \( H_{q,p}^{p,0} \) to denote the Fox’s H-function, defined by

(1.5)

\[ H_{q,p}^{p,0} \left( z \right) \left[ (β_q, A_q) \right] = \frac{1}{2π} \int_L \prod_{j=1}^p \Gamma(A_jz + α_j) \frac{ds}{s}, \]

where \( A_j, B_k > 0 \) and \( α_j, β_k > 0 \). The contour \( L \) has one of the following forms:

- \( L = L_{-∞} \) is a left loop in a horizontal strip starting at the point \( −∞ + iφ_1 \) and terminating at the point \( −∞ + iφ_2 \) with \( −∞ < φ_1 < φ_2 < ∞ \);
- \( L = L_{∞} \) is a right loop in a horizontal strip starting at the point \( ∞ + iφ_1 \) and terminating at the point \( ∞ + iφ_2 \) with \( ∞ < φ_1 < φ_2 < ∞ \);
- \( L = L_{iγ} \) is a contour starting at the point \( γ − i∞ \) and terminating at the point \( γ + i∞ + iφ_2 \), where \( γ ∈ R \).

Details regarding the contour and conditions for convergence of the integral in [17, 5] can be found in [16, Section 1.1.1.2] and [8].

In the course of our investigation, the first main tools is extended some results proved in [12] in the case when \( μ = −m, m ∈ N_0 \). Secondly, we establish the monotonicity of ratios involving the Fox-Wright functions.

2. Main results

The following theorem leads to an extension of the integral equation for the H-function obtained in [13] to the case \( μ = −m, m ∈ N_0 \).

Theorem 1. Suppose that \( μ = −m, m ∈ N_0 \) and \( \sum_{i=1}^p A_i = \sum_{j=1}^q B_j \). If \( γ ≥ 1 \), then the Fox-Wright function \( pΨq \) possesses the following integral representation

(2.6)

\[ pΨq \left[ (α_q, A_q) \right] = \int_0^1 e^{zt} H_{q,p}^{p,0} \left( t \frac{z}{(α_q, A_q)} \right) \frac{dt}{t} + \sum_{k=0}^m \sum_{j=0}^n \frac{t^{m−j}k^β_j}{k!}, \]

where the coefficients \( η \) and \( γ \) are defined by

(2.7)

\[ η = \left( 2π \right)^{p+1−q} \prod_{i=1}^p A_i^{α_q−1} \prod_{j=1}^q B_j^{β_q−1}, \quad γ = \min_{1 ≤ j ≤ p} (α_j/A_j), \]

and the coefficients \( l_r \) satisfy the recurrence relation:

(2.8)

\[ l_r = \frac{1}{r} \sum_{n=1}^r q_n l_r, \quad l_0 = 1, \quad \text{with} \quad q_n = \left( \frac{−1}{n} \right)^n \left[ \sum_{i=1}^p B_{n+1}(α_i)A_i − \sum_{j=1}^q B_{n+1}(β_j)B_j \right], \]

where \( B_n \) is the Bernoulli polynomial defined via generating function [3] p. 588

\[ \frac{te^{at}}{e^t − 1} = \sum_{n=0}^{∞} B_n(a) \frac{t^n}{n!}, \quad |t| < 2π. \]
Proof. In [1, Theorem 2], the authors found the Mellin transform of the $H$-function when $\mu = -m$, $m \in \mathbb{N}_0$, that is
\begin{equation}
\prod_{i=1}^p \Gamma(A_i k + \alpha_i) \prod_{j=1}^q \Gamma(B_j k + \beta_j) = \int_0^\rho H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right) t^{k-1} dt + \eta \rho \sum_{j=0}^m t_{m-j}^k, \Re(k) > \gamma.
\end{equation}
This implies that
\begin{align*}
p_{\Psi_q}^\left( (\alpha_p, A_p) \atop (\beta_q, B_q) \right) (z) &= \sum_{k=0}^\infty \prod_{i=1}^p \Gamma(A_i k + \alpha_i) \rho^{k} \\
&= \sum_{k=0}^\infty \int_0^\rho H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right) \left( \frac{(zt)^k}{k!} \right) dt + \sum_{k=0}^\infty \left( \eta \rho \sum_{j=0}^m t_{m-j}^k \right) \frac{k^j}{k!} \\
&= \sum_{k=0}^\infty \int_0^\rho H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right) \sum_{j=0}^m \left( \frac{(zt)^k}{k!} \right) dt + \eta \sum_{k=0}^\infty \sum_{j=0}^m \frac{t_{m-j}^k \rho^j}{k!} \\
&= \int_0^\rho e^{zt} H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right) dt + \eta \sum_{k=0}^\infty \sum_{j=0}^m \frac{t_{m-j}^k \rho^j}{k!}.
\end{align*}
For the exchange of the summation and integration, we use the asymptotic behavior of the $H$-function as $z \to 0 \ [\text{Theorem 1.2, Eq. 1.94}]
\begin{equation}
H_{q,p}^{m,n}(z) = \mathcal{O}(z^{-\gamma}), |z| \to 0,
\end{equation}
and the asymptotic behavior of the $H$-function as $z \to 0 \ [\text{Theorem 1}], we obtain
\begin{equation}
\int_0^\rho k^{j-1} \left| H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right. \right| dt \leq \int_0^\rho t^{-1} \left| H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right. \right| dt < M < \infty.
\end{equation}
Then, we are in position to apply the Lebesgue dominated convergence theorem. This completes the proof of Theorem [1].

Recall that a function $f : (0, \infty) \to (0, \infty)$ is called completely monotonic if $(-1)^n f^{(n)}(x) \geq 0$ for $x > 0$ and $n \in \mathbb{N}_0$. The celebrated Bernstein theorem asserts that completely monotonic functions are precisely those that can be expressed by the Laplace transform of a non-negative measure.

**Corollary 1.** Suppose that $\mu = 0$, $\gamma \geq 1$ and $\sum_{i=1}^p A_i = \sum_{j=1}^q B_j$. Then, the Fox-Wright function $p_{\Psi_q}^\left( (\alpha_p, A_p) \atop (\beta_q, B_q) \right)$ possesses the following integral representation
\begin{equation}
p_{\Psi_q}^\left( (\alpha_p, A_p) \atop (\beta_q, B_q) \right) (z) = \eta e^{-\rho z} = \int_0^\rho e^{-zt} H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right) dt, \ z \in \mathbb{R}.
\end{equation}
Moreover, if the function $H_{q,p}^{0,0}(\cdot)$ is non-negative, then the function
\begin{equation}
z \mapsto p_{\Psi_q}^\left( (\alpha_p, A_p) \atop (\beta_q, B_q) \right) (z) = \eta e^{-\rho z},
\end{equation}
is completely monotonic on $(0, \infty)$.

**Proof.** The application of Theorem [1] for $m = 0$ yields
\begin{equation}
p_{\Psi_q}^\left( (\alpha_p, A_p) \atop (\beta_q, B_q) \right) (z) = \int_0^\rho e^{zt} H_{q,p}^{0,0} \left( t \left| (B_q,\beta_q) \atop (A_p,\alpha_p) \right. \right) dt + \eta \rho z,
\end{equation}
which implies that (2.11) holds. Now, suppose that the $H$-function $H_{p,p}^{0,0}(\cdot)$, then by means of (2.11), we deduce that all prerequisites of the Bernstein Characterization Theorem for the complete monotone functions are fulfilled. 
\[\square\]
Example 1. The four parameters Wright function is defined by the series

\[ \phi((\mu_1, a), (\nu_1, b); z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(a + k\mu_1)\Gamma(b + k\nu_1)}, \quad \mu_1, \nu_1 \in \mathbb{R}, \quad a, b \in \mathbb{C}. \]  

The series on the right-hand side of (2.13) is absolutely convergent for all \( z \in \mathbb{C} \) if \( \mu_1 + \nu_1 > 0 \). If \( \mu_1 + \nu_1 = 0 \), the series is absolutely convergent for \( |z| < |\mu_1|^{1/|\nu_1|} \) and \( |z| = |\mu_1|^{1/|\nu_1|} \) under the condition \( \Re(a + b) > 2 \).

Some of the basic properties of the four parameters Wright function was proved in [4].

So, by means of formula (2.11) we deduce that the four parameters Wright function \( \phi((\mu_1, a), (\nu_1, b); z) \) admits the following integral representation:

\[ \phi((\mu_1, a), (\nu_1, b); z) = \int_0^{\mu_1 + \nu_1 - 1} e^{zt}H_{2,1}^{1,0}\left[ \begin{array}{c} (\mu_1, a), (\nu_1, b) \\ (1, 1) \end{array} \right] \frac{dt}{t} + \frac{\mu_1^{2-a} \nu_1^{2-b}}{\sqrt{2\pi}} e^{zt}t^{\nu_1}z^\nu, \]

where \( a, b, \mu_1, \nu_1 \in \mathbb{R} \) for which \( \mu_1 + \nu_1 = 1 \) and \( a + b = 3/2 \).

As a consequence, we derive the finite Laplace Transform for the function

\[ t \mapsto t^{-1}H_{2,1}^{1,0}\left[ \begin{array}{c} (1/2, 1/2), (1/2, 1) \\ (1, 1) \end{array} \right] \]

in \((0, 1/2)\). Recall that the finite Laplace Transform of a continuous (or an almost piecewise continuous) function \( f(t) \) in \((0, T)\) is denoted by

\[ \mathcal{L}_Tf(t) = \int_0^T e^{-st}f(t)dt. \]

We note that \( \mathcal{L}_Tf \) is actually the Laplace transform of the function \( f \) which vanishes outside of the interval \((0, T)\).

Example 2. Letting in (2.13) and (2.14), the values \( \nu_1 = \mu_1 = a = 1/2 \) and \( b = 1 \) and using the Legendre Duplication Formula

\[ \Gamma(z)\Gamma(z + 1/2) = 2^{1-2z}\sqrt{\pi}\Gamma(2z), \]

we get the following curious integral representation:

\[ \frac{e^{-2z} - e^{-z}}{\sqrt{\pi}} = \int_0^{1/2} e^{-zt}H_{2,1}^{1,0}\left[ \begin{array}{c} (1/2, 1/2), (1/2, 1) \\ (1, 1) \end{array} \right] \frac{dt}{t}, \quad z \in \mathbb{R}. \]

In the next result we show that the function \( p+1\Psi_q^\nu[\cdot] - \eta_1F_0[\cdot] \) is a Stieltjes transform.

Corollary 2. Let \( \sigma > 0 \) and \( z \in \mathbb{C} \), such that \( |\arg(1 + z)| < \pi \) and \( |z| < 1 \). In addition, assume that the hypotheses of Corollary[4] are satisfied. If \( H_{q,p}^{0,0}[\cdot] \) is non-negative, then, the following representation holds true:

\[ f(z) := \int_0^\rho H_{q,p}^{0,0}\left[ \begin{array}{c} (B_q, \beta_j) \\ (A_p, \alpha_p) \end{array} \right] - z - \eta \Gamma(\sigma, ; -; -\rho z) = \Gamma(\sigma) \int_0^\rho H_{q,p}^{0,0}\left[ \begin{array}{c} (B_q, \beta_j) \\ (A_p, \alpha_p) \end{array} \right] \frac{dt}{t(1 + tz)^\sigma}. \]

In particular, \( f(z) \) is completely monotonic on \((0, \infty)\).

Proof. Employing the generalized binomial expansion

\[ (1 + z)^{-\sigma} = \sum_{k=0}^\infty (\sigma)_k \frac{(-1)^k z^k}{k!}, \quad z \in \mathbb{C}, \quad |z| < 1, \]

with the formula (2.20) and the right hand side of (2.10) we obtain

\[ \Gamma(\sigma) \int_0^\rho H_{q,p}^{0,0}\left[ \begin{array}{c} (B_q, \beta_j) \\ (A_p, \alpha_p) \end{array} \right] \frac{dt}{t(1 + tz)^\sigma} = \Gamma(\sigma) \sum_{k=0}^\infty (\sigma)_k \frac{(-1)^k z^k}{k!} \left[ \int_0^\rho t^{k-1} H_{q,p}^{0,0}\left[ \begin{array}{c} (B_q, \beta_j) \\ (A_p, \alpha_p) \end{array} \right] \frac{dt}{t} \right] \]

\[ = \Gamma(\sigma) \sum_{k=0}^\infty (\sigma)_k \frac{(-1)^k z^k}{k!} \left[ \prod_{i=1}^\rho \frac{\Gamma(A_i k + \alpha_i)}{\Gamma(\beta_i + k j)} - \eta \rho^k \right] \]

\[ = \sum_{k=0}^\infty \Gamma(\sigma + k) \prod_{i=1}^\rho \frac{\Gamma(A_i k + \alpha_i)}{\Gamma(\beta_i + k j)} \frac{(-z)^k}{k!} - \eta \sum_{k=0}^\infty \Gamma(\sigma + k)(-\rho z)^k k! \]

\[ = p+1\Psi_q^\nu[\cdot] - z - \eta_1F_0[\cdot] - \eta \Gamma(\sigma, ; -; -\rho z). \]
Theorem 2. Suppose that \( \mu = -1 \) and \( \sum_{i=1}^{p} A_i = \sum_{j=1}^{q} B_j \). Then the following integral representation holds true:

\[
(2.20) \quad g(z) := p \Psi_{(\alpha_p, A_p)}^{(\beta_q, B_q)} \left[ \frac{z}{\eta (l_1 - \rho z)} \right] - \eta (l_1 - \rho z) e^{-\rho z} = \int_{0}^{\rho} e^{-zt} H_{p,p}^{0,0} \left( t \left| \left( \begin{array}{c} \beta_q, B_q \\ A_p, A_p \end{array} \right) \right. \right) dt, \quad z \in \mathbb{R}
\]

holds true. Moreover, if \( H_{p,p}^{0,0} \) is non-negative, then the function \( g(z) \) is completely monotonic on \((0, \infty)\).

Proof. From Theorem 1, when \( \mu = -1 \), we have

\[
(2.19) \quad p \Psi_{(\alpha_p, A_p)}^{(\beta_q, B_q)} \left[ \frac{z}{\eta (l_1 - \rho z)} \right] - \eta (l_1 - \rho z) e^{-\rho z}.
\]

\[
\square
\]

Example 3. The four parameters Wright function \( \phi ((\mu, a), (\nu, b); z) \) possesses the following integral representation

\[
(2.20) \quad \phi ((\mu, a), (\nu, b); z) = \int_{0}^{\mu + \nu} e^{-zt} H_{2,1}^{0,1} \left[ \left( \begin{array}{c} \mu, a \\ \nu, b \end{array} \right) \right] dt.
\]

where

\[
(2.20) \quad \phi ((\mu, a), (\nu, b); z) = \int_{0}^{\mu + \nu} e^{-zt} H_{2,1}^{0,1} \left[ \left( \begin{array}{c} \mu, a \\ \nu, b \end{array} \right) \right] dt.
\]

The following lemma is called the Jensen’s integral inequality, for more details, one may see [5, Chap. I, Eq. (7.15)].

Lemma 1. Let \( \omega \) be a non-negative measure and let \( \varphi \geq 0 \) be a convex function. Then for all \( f \) be a integrable function we have

\[
(2.21) \quad \varphi \left( \int f d\nu \right) \leq \int \varphi(f) d\nu.
\]

In the next theorem we present a new Luke type inequality when \( \mu = 0 \).

Theorem 2. Keep the notations and constraints of hypotheses of Corollary 7. Assume that the function \( H_{q,r}^{m,n} \) is non-negative, then the following two-sided bounding inequality holds true:

\[
(2.22) \quad \Psi_0 e^{-\Psi_0 z} + \eta e^{-\rho z} \leq p \Psi_{q}^{(\alpha_p, A_p)} \left[ \frac{z}{\eta} \right] \leq \left( \Psi_0 - \frac{\Psi_1}{\rho} \right) + \eta e^{-\rho z},
\]

where

\[
\Psi_0 := \prod_{i=1}^{p} \Gamma(\alpha_i) \prod_{j=1}^{q} \Gamma(\beta_j) - \eta,
\]

\[
\Psi_1 := \prod_{i=1}^{p} \Gamma(\alpha_i + A_i) \prod_{j=1}^{q} \Gamma(\beta_j + B_j) - \eta \rho.
\]

Proof. Letting \( \varphi(t) = e^{-zt}, \ f(t) = t \), and

\[
(2.20) \quad d\nu(t) = H_{p,p}^{0,0} \left( t \left| \left( \begin{array}{c} A_p \\ A_p \end{array} \right) \right. \right) dt.
\]

From (2.20) we get

\[
\int_{0}^{\rho} d\nu(t) = \prod_{i=1}^{p} \Gamma(\alpha_i) \prod_{j=1}^{q} \Gamma(\beta_j) - \eta = \Psi_0, \quad \text{and} \quad \int_{0}^{\rho} f(t) d\nu(t) = \prod_{i=1}^{p} \Gamma(\alpha_i + A_i) \prod_{j=1}^{q} \Gamma(\beta_j + B_j) - \eta \rho = \Psi_1.
\]

and using (2.20) when \( m = 0 \) we find

\[
\int_{0}^{\rho} \varphi(f(t)) d\nu(t) = p \Psi_{q}^{(\alpha_p, A_p)} \left[ \frac{z}{\eta} \right] - \eta e^{-\rho z}.
\]
Hence, Lemma 4 completes the proof of the lower bound of inequalities (2.22). In order to demonstrate the upper bound, we will apply the converse Jensen inequality, due to Lah and Ribarić, which reads as follows. Set

\[ A(f) = \int_m^M f(s) \, d\sigma(s) / \int_m^M \, d\sigma(s), \]

where \( \sigma \) is a non-negative measure and \( f \) is a continuous function. If \(-\infty < m < M < \infty\) and \( \varphi \) is convex on \([m, M]\), then according to [15, Theorem 3.37]

\[ (M - m)A(\varphi(f)) \leq (M - A(f)) \varphi(m) + (A(f) - m) \varphi(M). \]

Setting \( \varphi(t) = e^{-zt} \), \( d\sigma(t) = dv(t) \), \( f(s) = s \) and \([m, M] = [0, \rho]\), we complete the proof of the upper bound in (2.22).

In view of inequalities (2.22) and the Laplace transform of the function \( x^{\lambda-1} \psi(x) \) [18, Eq. (7)]

\[ \int_0^\infty e^{-t\lambda-1} \psi(x) \left[ (\alpha_p, \alpha_p) \right] zt \, dt = p+1 \psi(x) \left[ (\lambda, \alpha_p) \right] z, \]

and make use of the following known formula

\[ \int_0^\infty t^{\lambda-1} e^{-\sigma t} \, dt = \frac{\Gamma(\lambda + 1)}{\sigma^{\lambda+1}}, \quad (\lambda > -1, \, \sigma > 0), \]

we can deduce the new following inequalities for the function \( p+1 \psi(x) \) :

**Corollary 4.** Let \( \lambda > 0 \) and suppose the hypotheses of Corollary 7 are satisfied. If the function \( H_{q,\rho}^{0,p}[.] \) is non-negative, then the following two-sided bounding inequality holds true:

\[ \eta \frac{\Gamma(\lambda)}{(1 + \rho z)^\lambda} + \psi(0) \frac{\Gamma(\lambda)}{(1 + (\psi/\psi_0)z)^\lambda} \leq p+1 \psi(x) \left[ (\lambda, \alpha_p, \alpha_p) \right] z \leq \frac{\Gamma(\lambda) \psi(0)}{p} + \frac{\Gamma(\lambda) \left( \eta + \frac{\rho z}{\psi} \right)}{(1 + \rho z)^\lambda}. \]

**Theorem 3.** Keep the notations and constraints of hypotheses of Corollary 7. Assume that the function \( H_{q,\rho}^{0,p}[.] \) is non-negative, then the following inequality

\[ \frac{\Gamma(\sigma) \psi(0)}{(1 + (\psi/\psi_0)z)^\sigma} + \eta 1 F_0(\sigma; \sigma; -\rho z) \leq p+1 \psi(x) \left[ (\lambda, \alpha_p, \alpha_p) \right] z, \]

is valid for all \( \sigma > 0 \) and \( |z| < 1 \).

**Proof.** We set \( \varphi(u) = u^\sigma, \, \sigma > 0, \, f(t) = 1/(1 + tz) \) and \( dv(t) = t^{-1} \Gamma(\sigma) H_{q,\rho}^{0,p} \left[ l^{(B_\beta, \beta)} \right] dt \). By (2.22) we have

\[ \int_0^\rho dv(t) = \Gamma(\sigma) \int_0^{\rho \Gamma(\sigma)} H_{q,\rho}^{0,p} \left[ l^{(B_\beta, \beta)} \right] dt \]

\[ = \Gamma(\sigma) \left( \prod_{j=1}^{\rho} \Gamma(\alpha_j) \right) - \eta \Gamma(\sigma). \]

Moreover (2.16), reads

\[ \int_0^\rho f(t) dv(t) = p+1 \psi(x) \left[ (\lambda, \alpha_p, \alpha_p) \right] z - \frac{\eta}{1 + \rho z}. \]

and

\[ \int_0^\rho \varphi(f(t)) dv(t) = p+1 \psi(x) \left[ (\lambda, \alpha_p, \alpha_p) \right] z - \frac{\eta}{1 + \rho z}. \]

By means of Lemma 1 we obtain

\[ \Gamma(\sigma) \psi(0)^{1-\sigma} \left( p+1 \psi(x) \left[ (\lambda, \alpha_p, \alpha_p) \right] z - \frac{\eta}{1 + \rho z} \right)^{\sigma} \leq p+1 \psi(x) \left[ (\lambda, \alpha_p, \alpha_p) \right] z - \eta 1 F_0(\sigma; \sigma; -\rho z). \]

By virtue of the left-hand side of inequality (2.22) and (2.30), we conclude the inequality (2.24). \qed

The next lemma is in fact the so-called the Chebyshev integral inequality, see [5, p. 40].
Lemma 2. If \( f, g : [a, b] \to \mathbb{R} \) are synchronous (both increasing or decreasing) integrable functions, and \( p : [a, b] \to \mathbb{R} \) is a positive integrable function, then

\[
\int_a^b p(t)f(t)dt \int_a^b p(t)g(t)dt \leq \int_a^b p(t)dt \int_a^b p(t)f(t)g(t)dt.
\]

Note that if \( f \) and \( g \) are asynchronous (one is decreasing and the other is increasing), then (2.31) is reversed.

Theorem 4. Assume that the hypotheses of Corollary 4 are satisfied. Suppose that \( \delta, \sigma > 0 \). If \( H_{p,q}^{0,0} \) is non-negative, then the function

\[
F := F_{(\sigma,1),\rho}^{(\sigma,1),\rho}(z) = \frac{p+1}{p} \left( \int_0^p H_{p,q}^{0,0}\left[ (A_{\sigma,\rho}, A_{\sigma,\rho}) \right] \left[ \frac{(1+z)^{\rho \sigma}}{(1+z)^{\rho \sigma + 1}} \right] \right)
\]

is increasing on \((0,1)\). In addition, the function \( F(z) \) is decreasing on \((0,1)\) for each \( \delta < 0 \) and \( \sigma > 0 \).

Proof. In view of (2.16), using the following property of the Fox \( H \)-function [3] Property 1.5, p. 12

\[
H_{q,p}^{n,m}\left[ (A_{\sigma,\rho}, A_{\sigma,\rho}) \right] = z^\delta H_{q,p}^{n,m}\left[ (A_{\sigma,\rho}) \right],
\]

we can rewrite the function \( F \) as follows:

\[
F_{(\sigma,1),\rho}^{(\sigma,1),\rho}(z) = \frac{p+1}{p} \left( \int_0^p H_{p,q}^{0,0}\left[ (A_{\sigma,\rho}, A_{\sigma,\rho}) \right] \frac{dt}{(1+z)^{\rho \sigma + 1}} \right)
\]

Now, we consider the functions \( p, f, g : [0, \rho] \to \mathbb{R} \), defined by

\[
p(t) = t^{-1}(1+z)^{-\sigma} H_{p,q}^{0,0}\left[ (A_{\sigma,\rho}) \right], f(t) = t^\sigma, g(t) = \frac{t}{1+z}.
\]

Observe that the functions \( f \) and \( g \) are increasing, thus, by means of Lemma 2 we infer

\[
\int_0^p \frac{dt}{t} \left( \int_0^p H_{p,q}^{0,0}\left[ (A_{\sigma,\rho}, A_{\sigma,\rho}) \right] \frac{dt}{(1+z)^{\rho \sigma + 1}} \right)
\]

On the other hand, we have

\[
\frac{1}{\sigma} \left( \int_0^p H_{p,q}^{0,0}\left[ (A_{\sigma,\rho}, A_{\sigma,\rho}) \right] \frac{dt}{(1+z)^{\rho \sigma + 1}} \right)^2 \frac{\partial}{\partial z} F_{(\sigma,1),\rho}^{(\sigma,1),\rho}(z) =
\]

By (2.33) and (2.34) we deduce that the function \( z \to F(z) \) is increasing on \((0,1)\) for all \( \sigma > 0 \) and \( \delta > 0 \). Moreover, if \( \delta < 0 \) then the inequality (2.34) is reversed and consequently the function \( z \to F(z) \) is decreasing on \((0,1)\) for all \( \sigma > 0 \) and \( \delta < 0 \). Now, the proof of this theorem is completed.

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