CIRCUMBINARY RING, CIRCUMSTELLAR DISKS, AND ACCRETION IN THE BINARY SYSTEM UY AURIGAE

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ABSTRACT

Recent exo-planetary surveys reveal that planets can orbit and survive around binary stars. This suggests that some fraction of young binary systems which possess massive circumbinary (CB) disks may be in the midst of planet formation. However, there are very few CB disks detected. We revisit one of the known CB disks, the UY Aurigae system, and probe 12CO 2–1, 13CO 2–1, C18O 2–1, SO 5(6)–4(5) and 12CO 3–2 line emission and the thermal dust continuum. Our new results confirm the existence of the CB disk. In addition, the circumstellar (CS) disks are clearly resolved in dust continuum at 1.4 mm. The spectral indices between the wavelengths of 0.85 mm and 6 cm are found to be surprisingly low, being 1.6 for both CS disks. The deprojected separation of the binary is 1.26 based on our 1.4 mm continuum data. This is 0.07 (10 AU) larger than in earlier studies. Combining the fact of the variation of UY Aur B in R band, we propose that the CS disk of an undetected companion UY Aur Bb obscures UY Aur Ba. A very complex kinematical pattern inside the CB disk is observed due to a mixing of Keplerian rotation of the CB disk, the infall and outflow gas. The streaming gas accreting from the CB ring toward the CS disks and possible outflows are also identified and resolved. The SO emission is found to be at the bases of the streaming shocks. Our results suggest that the UY Aur system is undergoing an active accretion phase from the CB disk to the CS disks. The UY Aur B might also be a binary system, making the UY Aur a triple system.

Key words: planet–disk interactions – protoplanetary disks – stars: formation – stars: individual (UY Aurigae)

Online-only material: color figures

1. INTRODUCTION

A substantial fraction of stars form in binary or multiple systems (Duchêne & Kraus 2013) and many recent results of exo-planetary surveys reveal that planets can orbit and survive both around the individual components of a binary star and in circumbinary (CB) orbits (Udry & Santos 2007; Roell et al. 2012). This suggests that at least a reasonable fraction of young binary systems exhibit circumstellar (CS) disks and CB rings. To be gravitationally stable, the CS disks are expected to be located inside the Roche lobe and the CB rings are expected to be outside the outer Lindblad resonances (or the 4:1 resonance for a binary with low to moderate eccentricity; Artymowicz et al. 1991). Any residual material inside the non-stable intermediate area in between the CB rings and CS disks would inflow toward the CS disks through streamers. These streamers feed the CS disks, which would otherwise dissipate within a relatively short timescale. Signs of accretion and near-infrared (NIR) excesses are observed around many T Tauri binary systems of a few million years. Thus, streamers are expected to be common and to provide enough material from the outer CB disk to maintain the inner CS disks, to sustain accretion onto stars, and (maybe) to form planets. Only a few sources have been observed with clear signatures of streamers. The best evidence in binary systems remains the infrared image of SR 24 (Mayama et al. 2010), although the kinematics and mass of the infalling material were not directly measured.

With the exception of the GG Tau ring (Dutrey et al. 1994) and of the CB disk around V4046 Sgr (Rosenfeld et al. 2013), the number of CB disks detected so far is also surprisingly small, partly because of the limited sensitivity of millimeter/submillimeter arrays prior to ALMA. The formation and survival of a fully shaped CB disk requires very specific physical conditions. (Magneto-) hydrodynamic ((M)HD) simulations show that this depends on the initial angular momentum, the binary mass ratio (Bate & Bonnell 1997), and the strength and distribution of the magnetic field in the initial molecular core (Machida et al. 2008). In fact, spiral arms form and they may be less mimetic a ring shape (for illustration, see Figure 5 from Machida et al. 2008 or Figure 2 from Bate & Bonnell 1997). This is particularly true if the brightness contrast between the main spiral arm and the others is significant (10 or more) under moderate angular resolutions with respect to the width of the main spiral arm. Thanks to the resolving power and high contrast of new NIR facilities, spiral patterns have recently been resolved around many young systems (Fukagawa et al. 2010). They are also seen in recent images obtained from millimeter/submillimeter arrays such as those of the disks orbiting MW758 (Isella et al. 2010) or AB Aurigae (Tang et al. 2012). In most cases, more observations (with an adequate hydrodynamical modeling) are needed to fully characterize the behavior and origin of such dynamical disturbances.

To study the origin and evolution of accretion processes in young multiple T Tauri systems, we have chosen to focus
our interest on UY Aurigae (hereafter, UY Aur). UY Aur is among the few binary systems that exhibit a ring-like CB disk (Close et al. 1998). It is a system separated by \(0\,^{\prime\prime}88\) in the Taurus–Aurigae star-forming region. We assume UY Aur is at a distance of 140 pc. Based on the spectral energy distribution, the UY Aur stellar components are classified as class II objects (Beckwith et al. 1990). The primary, UY Aur A, has a spectral type K7 (0.95 \(M_\odot\)) and the secondary, UY Aur B, has a spectral type M0 (0.6 \(M_\odot\)) based on the photometric absorption features (Herbst et al. 1995). A more recent estimate of the binary properties, based on the spectroscopy survey by Hartigan & Kenyon (2003), provides a spectral type of M0 (0.6 \(M_\odot\)) and M2.5 (0.34 \(M_\odot\)) for UY Aur A and B, respectively.

The CB and CS disks of UY Aur have been observed at various wavelengths. At optical/NIR, a bright and patchy dusty CB ring inclined by \(\sim-42^\circ\) and CS materials connecting the CB ring and the CS disks were reported from observations with Canada–France–Hawaii Telescope (Close et al. 1998) and with SUBARU (see Figure 1; Hioki et al. 2007). The brightest spot at NIR is on the south CB ring called NIR 1 (Hioki et al. 2007). The CB disk itself is detected in \(^{13}\)CO 2–1 and 2–1 and yields an enclosed mass of about 1.2 \(M_\odot\) (Duvert et al. 1998). Besides the CB disk, there are extended ridges and hint of streamers seen in the measurements of the CO isotopologues. At 3.6 cm, the binary has also been detected using the Very Large Array (VLA) with a 2.3 beam (Contreras & Wilkin 2006), which does not separate the two sources. The spectral index of the UY Aur system was found to be surprisingly low, being \(\sim-1.6\) between 3.6 cm and 1.4 mm (Contreras & Wilkin 2006).

The photometric properties of the UY Aur binary are peculiar. Until recently, UY Aur B has a high extinction at NIR while still exhibiting a stellar spectrum (Herbst et al. 1995). The \(R\)-band magnitude of UY Aur B changed by \(\sim5\) mag within 50 yr. The \(H\)-band (\(\lambda = 1.6 \mu m\)) magnitude of UY Aur B varied by 1.3 mag. UY Aur A, which is relatively stable in the \(H\) band (\(\Delta H = 0.4\) mag) between 1996 and 2005 (Close et al. 1998; Hioki et al. 2007), exhibits a strong optical variability in the \(V\) and \(B\) bands (Berdnikov et al. 2010).

In this paper, we report our new observational results of the \(^{13}\)CO 2–1, \(^{18}\)O 2–1, SO 5(6)–4(5), and \(^{12}\)CO 3–2 molecular transitions and of the thermal dust emissions at 1 mm and 3 mm.

2. OBSERVATIONS

2.1. IRAM Plateau de Bure Interferometer

The observations were carried out in nine tracks on the IRAM Plateau de Bure Interferometer (PdBI) from 2011 December 26 to 2012 March 16. At 1.4 mm (220 GHz), the observations were carried out in the B, C, and D configurations, covering the baseline range of 15 m (11 \(k\alpha\)) to 456 m (338 \(k\alpha\)), in dual linear polarization. The broadband correlator provided an effective bandwidth of 7.6 GHz. The total on source time is 44 hr. The narrow-band correlator was set to cover the \(^{13}\)CO \(J = 2–1\) and \(^{18}\)O \(J = 2–1\) lines with a spectral resolution of 39 kHz (0.053 km s\(^{-1}\)). The SO 5(6)–6(5) is detected in the broadband correlator with a spectral resolution of 1.953 MHz (2.662 km s\(^{-1}\)). At 3 mm (100 GHz), the observation were carried out in the A configuration, covering the baseline range of 110.5 m (37 \(k\alpha\)) to 666 m (223 \(k\alpha\)), with an on source time of 6 hr. All data were processed using the CLIC software in the GILDAS package. Phase, amplitude, bandpass, and flux calibration were performed following the standard procedure using 0400+258, 0459+252 for the amplitude and phase, 3C84 for the bandpass, and MWC349 for the flux. To improve the image quality, we have applied phase-only self-calibration with a time interval of 45 s on the 1.4 mm continuum, and transfer the solution to the spectral line data. For display, the channel maps of the \(^{13}\)CO 2–1 and \(^{18}\)O 2–1 were smoothed in velocity by three-channel averaging.

The phase center is \(\alpha = 04:51:47.388, \delta = 30:47:13.1\). The effective angular resolution varies depending on the weighting scheme during inverse Fourier transformation to the image domain. The highest angular resolution of 0.55 \(\times\) 0.34 for the 1.4 mm continuum image is achieved using uniform weighting in order to resolve UY Aur A and B, with an effective noise level (dynamic range limited) of 0.18 mJy beam\(^{-1}\) or 1.8 mK in brightness. At 3 mm, we reach an angular resolution of 0.97 \(\times\) 0.54 using robust weighting parameter 1 and a noise level of 0.04 mJy beam\(^{-1}\), or 9.6 mK. For \(^{13}\)CO, the angular resolution is 1.27 \(\times\) 0.92, obtained with robust weighting parameter 1 and the brightness sensitivity is 17 mJy beam\(^{-1}\), or 0.36 K with a channel spacing of 0.053 km s\(^{-1}\) (the spectral resolution being 1.6 times coarser due to channel shape). The \(^{18}\)O maps presented in Figure 5 have been tapered to a lower spectral angular resolution and the natural weighting scheme is used. The angular resolution is 1.94 \(\times\) 1.74 and the noise level is 12 mJy beam\(^{-1}\), or 0.09 K. The SO maps are obtained with robust weighting parameter 1 with an angular resolution of 1.25 \(\times\) 0.99 and a noise level of 1.2 mJy beam\(^{-1}\).

2.2. The Submillimeter Array

In addition to the PdBI observations, we also obtained the \(^{12}\)CO 3–2 line data from the Submillimeter Array (SMA).\(^8\) The observations were carried out on 2013 November 24 in the extended configuration. The newly obtained data are combined with archival data taken in compact configuration on 2010 December 13 in order to have better \(uv\) coverage. The data were reduced using the MIR package following the standard calibration process. The spectral resolution is 0.35 km s\(^{-1}\). The \(^{12}\)CO 3–2 maps are made with natural weighting, giving an angular resolution of 1.12 \(\times\) 0.98 at position angle (P.A.) 123\(^\circ\). The noise level at 1\(\sigma\) is 0.17 Jy beam\(^{-1}\), or 1.6 K.

3. RESULTS

3.1. Continuum

3.1.1. Compact Continuum Emission

Continuum emission is detected at all three frequencies. At 1.4 mm, the emission is detected toward UY Aur A and B, which are clearly separated with the 0.4 beam. The 1.4 mm and 3 mm continuum overlaid on a SUBARU image are shown in Figure 1. The coordinates and flux densities of A and B are determined by fitting the visibilities with two Gaussians. The best-fit J2000 coordinates at the epoch of observations are (\(\alpha, \delta\)) = (04:51:47.392, 30:47:13.22) and (\(\alpha, \delta\)) = (04:51:47.335, 30:47:12.63) for the continuum emission from UY Aur A and B, respectively. The flux densities of the UY Aur sources between 0.85 mm and 6.0 cm are listed in Table 1. We note that the flux density at 0.85 mm quoted in Table 1 are from Akeson & Jensen (2014) due to a much higher angular resolution as compared to our SMA 0.85 mm measurement. As

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\(^8\) The Submillimeter Array is a joint project between the Smithsonian Astrophysical Observatory and the Academia Sinica Institute of Astronomy and Astrophysics and is funded by the Smithsonian Institution and the Academia Sinica.
Figure 1. (a) Continuum emission at 1.4 mm in white contours at levels 1, 4, 7, 15, and 25 × 0.54 (3σ) mJy per 0′′/55 × 0′′/34 with position angle (P.A.) 24′. beam. (b) Continuum emission at 3 mm in black contours at 1, 5, 10, 20, 30, and 40 × 0.11 (3σ) mJy per 0′′/97 × 0′′/54 beam with P.A. 45′. (c) Residual continuum emission after removing two Gaussian fits at UY Aur A and B. Contour steps are 3σ (0.323 mJy per 1′′/9 × 1′′/7 beam). In all panels, the H-band coronagraphic image by SUBARU (Hioki et al. 2007) is shown in color scale. The white ellipse marks the 2.1 μm peak intensity of the disk with a semi-major axis of 3′/7 and inclination of 42′ determined by Close et al. (1998). Note that the ellipse center is 0′.7 off in Declination from the phase center. The white square marks the peak at 2.1 μm and labeled as NIR 1 component by Hioki et al. (2007). The gray triangle marks the SO southeast peak (see Section 3.2.2). The crosses mark the locations of UY Aur A and B determined using our 1.3 mm continuum image at the epoch of observation.

(A color version of this figure is available in the online journal.)

Figure 2. Spectral energy distribution plots of UY Aur A and B. The triangle marks the 3σ upper limit of $S_A$ and $S_B$ at 6 cm.

### Table 1

| λ (mm) | $S_A$ (mJy) | $S_B$ (mJy) | Reference |
|--------|-------------|-------------|-----------|
| 0.85   | 48.40 ± 0.66 | 18.27 ± 0.72 | ALMA; Akeson & Jensen (2014) |
| 1.4    | 22.78 ± 0.09 | 6.81 ± 0.09  | PdBI; this work |
| 3.4    | 5.4 ± 0.1    | 1.7 ± 0.1    | PdBI; this work |
| 36     | 0.12 ± 0.03  | 0.11 ± 0.03  | VLA; Contreras & Wilkin (2006) |
| 60     | <0.12        | <0.12        | VLA; data archive |

**Notes.** Columns are wavelengths (λ) in millimeters, flux density in mJy of UY Aur A ($S_A$) and UY Aur B ($S_B$), and the references.

Discussed in Contreras & Wilkin (2006), the 3.6 cm emission is consistent with the stellar location of the binary. At 6 cm, we obtained the archival data of the VLA and analyzed them. There is no emission detected within the field of view, suggesting an upper limit of 0.12 mJy at 3σ. We derive a spectral index $\alpha$, defined as $S(\nu) = S(\nu_0)(\nu/\nu_0)^\alpha$, below 2 for both sources ($\approx$1.6; see Figure 2).

At 1.4 mm, the emission around UY Aur A is well represented by a circular Gaussian of deconvolved size $\sim$0.2. The deconvolved size of $\sim$0.2 probably reflects the residual seeing after self-calibration. On the other hand, the emission around UY Aur B appears more extended: an elliptical Gaussian fit yields a major axis ranging from 0′′34 to 0′′45, oriented along the AB line. Using uniform disk models would yield larger sizes of the disks, i.e., 0′′34 and 0′′47 for A and B, respectively. At 3 mm, UY Aur B also appears slightly more extended (0′′17) than UY Aur A (0′′08), although we have to fix the relative positions because of the more limited angular resolution.

Assuming the whole system is coplanar with the CB ring, we deproject the visibilities of the 1.4 mm continuum data with an inclination angle ($i$) of 42° along P.A. 135° in order to ease the comparison with previous studies. The emission from UY Aur A again appears compatible with a circular Gaussian of size of 0′′22 (30 AU), which, as mentioned in the previous paragraph, is most likely an upper limit due to seeing limitations. The 1.4 mm emission around UY Aur B is fitted by an elliptical Gaussian with a best-fit major axis of 0′′46 (65 AU) at P.A. 243° ± 7°, i.e., 18° ± 7° away from the rotation axis of the system. The fitting results are listed in Table 2.

We also derive a (deprojected) separation between the millimeter emissions around A and B of 1′′26 ± 0′′013 at P.A. = 229°6 ± 0°3. This separation can be compared to the optical measurements reported over the last decades (see Table 3), which are compatible with a circular orbit of radius $\sim$1′′19 (167 AU Hioki et al. 2007). The millimeter separation is thus larger by 0′′07 at a 5σ level. If we allow the orbit to be elliptical by fitting a linear slope in the deprojected separation, the best-fit slope is $1.2 \pm 0.4$ mas yr$^{-1}$ (Figure 3) and the 1.4 mm position still deviates by $\sim$0′′05 (4σ) from the extrapolated optical position. We note, however, that the inclination angle is uncertain. A more detailed assessment of the UY Aur system inclinations will be presented in Section 4.4.

### 3.1.2. Extended Continuum Emission

At high spatial resolution, there is no visible emission outside the stellar positions, in particular around the CB
disk/ring traced by NIR scattered light. However, the rms noise level, 0.18 mJy beam$^{-1}$, is relatively high at this resolution. To trace the fainter extended continuum emission, we reconstructed the residual image after subtracting the best-fit Gaussian of UY Aur A and B. We used natural weighting with a corresponding resolution of 1.9 $\times$ 1.69 in order to increase the brightness sensitivity, reaching 0.11 mJy beam$^{-1}$ or 0.86 mK. The total flux in this residual image is 7.6 mJy. The total flux emanating from UY Aur system at 1.4 mm is thus about 38 mJy. This matches well with the 48 $\pm$ 9 mJy measured at 1.2 mm by Beckwith et al. (1990), which gives 37 $\pm$ 6 mJy after correction for the mean spectral index of the emission. We have thus recovered essentially all the emission.

The resulting image is overlaid on the NIR emission in Figure 1. The 1.4 mm residual emission is brightest at 1$^\circ$5 away from the binary, especially to the north and to the east, while the NIR emission is brighter south and west. There is substantial millimeter emission near the northwest CS material identified by Hioki et al. (2007) or the possible spiral arm identified by Close et al. (1998), hereafter NW spiral, connecting from the CB ring to the inner disks. Assuming a temperature of 20 K, and an opacity of gas and dust of 0.02 cm$^2$ g$^{-1}$, the upper limit (at 3$\sigma$ = 2.6 mK) of the H$_2$ column density, $N_{\text{CB, gas}}$, is 4.3 $\times$ 10$^{22}$ cm$^{-2}$ in the CB ring. The detected extended continuum peaks have column densities two to four times higher. The total continuum flux (7.6 mJy) recovered in the CB ring yields a total mass of $\sim$7 $\times$ 10$^{-4} M_\odot$.

### 3.2. Gas

#### 3.2.1. The Extended Gas Traced with CO 2$\rightarrow$1 Isotopologues

The $^{13}$CO 2$\rightarrow$1 and C$^{18}$O 2$\rightarrow$1 emissions are well detected and resolved by the new observations. The channel maps of $^{13}$CO and C$^{18}$O are shown Figures 4 and 5, respectively. The total intensity maps of the $^{13}$CO 2$\rightarrow$1, C$^{18}$O 2$\rightarrow$1, SO 5(6)$\rightarrow$4(5), and 12CO 3$\rightarrow$2 are presented in Figure 7. At low velocities, most of the extended emissions trace the CB disk/ring with peaks along the northwest and southeast directions, consistent with the proposed major axis of 135$^\circ$ from the NIR image and with the previous images by Duvert et al. (1998). However, the new millimeter data have twice the higher angular resolution.

### Table 2

Continuum Best-fit Deconvolved Size

| Fitting Function | $\theta_{\text{maj}}, \theta_{\text{min}}, \text{P.A.}$ ("), ("), (°) | Fitting Function | $\theta_{\text{maj}}, \theta_{\text{min}}, \text{P.A.}$ ("), ("), (°) | $\lambda$ (mm) |
|-----------------|---------------------------------|-----------------|---------------------------------|----------------|
| Circular-Gauss  | 0.21 $\pm$ 0.01                | Circular-Gauss  | 0.31 $\pm$ 0.02                | 1.4            |
| Circular-Gauss  | 0.21 $\pm$ 0.01                | Elliptical-Gauss| 0.45 $\pm$ 0.03, [0.2], 57 $\pm$ 5 | 1.4            |
| Elliptical-Gauss| 0.24 $\pm$ 0.01,0.19 $\pm$ 0.01,110 $\pm$ 7 | Elliptical-Gauss| 0.34 $\pm$ 0.03, [0.2], 55 $\pm$ 6 | 1.4            |
| Circular-Gauss  | 0.08 $\pm$ 0.03                | Circular-Gauss  | 0.17 $\pm$ 0.05                | 3.4            |

#### Notes.

Columns are the fitting function (f), the de-convolved size ($\theta$ in arcsecond) of A and B, and the observed wavelengths ($\lambda$ in millimeters). The major and minor diameters and the P.A. are given. The deprojection of the global system is done with an inclination of 42$^\circ$ and P.A. of 135$^\circ$. The minor axis of the elliptical Gaussian fit toward UY Aur B is fixed to 0." $\times$ 1"

### Table 3

Astrometry Parameters

| Date           | P.A.   | $q$   | P.A.$\text{deproj.}$ | $q_{\text{deproj.}}$ | Reference           |
|----------------|--------|-------|----------------------|----------------------|---------------------|
| 1944.25        | 212.2  | 0.82  | 215.4                | 1.09                 | Joy & van Biesbroeck (1944) |
| 1990 Nov 10    | 227.1  | 0.875 | 226.6                | 1.1771               | Ghez et al. (1995)   |
| 1993 Dec 26    | 227.6  | 0.861 | 226.9                | 1.158                | Ghez et al. (1995)   |
| 1994 Sep 22    | 226.5  | 0.868 | 226.1                | 1.168                | Ghez et al. (1995)   |
| 1994 Oct 18    | 228.2  | 0.887 | 227.4                | 1.193                | Ghez et al. (1995)   |
| 1996 Oct 24    | 228.1  | 0.88  | 227.1                | 1.18                 | Close et al. (1999)  |
| 1997 Dec 8     | 227.1  | 0.878 | 226.6                | 1.181                | White & Ghez (2001)  |
| 1998.98        | 221.7  | 0.892 | 227.3                | 1.200                | Hartigan & Kenyon (2003) |
| 1999 Nov 16    | 227.6  | 0.88  | 226.9                | 1.18                 | McCabe et al. (2006) |
| 2000 Feb 22–23 | 228.8  | 0.894 | 227.84                | 1.202                | Brandeker et al. (2003) |
| 2002 Nov 23    | 229.13 | 0.886 | 228.07                | 1.191                | Hioki et al. (2007)  |
| 2005 Nov 9     | 231.36 | 0.890 | 229.74                | 1.194                | Hioki et al. (2007)  |
| 2011 Dec 30    | 231.22 | 0.94  | 229.63                | 1.26                 | This work            |

#### Notes.

Columns are the observation date (date), the P.A. of the orientation between A and B (P.A.), the observed separation (q), the de-projected P.A. of the orientation (P.A.$\text{deproj.}$), the de-projected separation $q_{\text{deproj.}}$, and the references. The deprojection is calculated assuming i = 42$^\circ$ and P.A. of the CB disk of 45$^\circ$. 

The total flux in this residual image is 7.6 mJy. The total flux emanating from UY Aur system at 1.4 mm is thus about 38 mJy. This matches well with the 48 $\pm$ 9 mJy measured at 1.2 mm by Beckwith et al. (1990), which gives 37 $\pm$ 6 mJy after correction for the mean spectral index of the emission. We have thus recovered essentially all the emission.

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Section 3.3.3), we estimate the gas column density, \(N_{\text{H}_2}\), of the CB ring from the \(^{13}\text{CO} 2–1\) line. The total emission of the CB ring except at NIR 1, which is probably contaminated by another cloud, is \(\sim 0.49\) Jy beam\(^{-1}\) km s\(^{-1}\) (or 10.4 K km s\(^{-1}\)). We assume an excitation temperature of \(\sim 20\) K, an isotopic ratio \(^{13}\text{C}/^{12}\text{C}\) of 77 (Wilson & Rood 1994, the value in local interstellar medium, ISM), and a CO/H\(_2\) ratio of \(10^{-4}\). We find \(N_{\text{H}_2} = 4 \times 10^{21}\) cm\(^{-2}\), consistent with the upper limit of \(4.3 \times 10^{22}\) cm\(^{-2}\) (Section 3.1) based on the 1.4 mm continuum. Thus, the detection of the CB ring dust continuum would require an improvement in sensitivity by about one order of magnitude (if the dust absorption coefficient per unit mass is correct), down to 8 \(\mu\)Jy beam\(^{-1}\) at 1.4 mm. The total gas mass (including both the gas mass of the ring and of the filaments) is \(M_{\text{gas}} \sim 2 \times 10^{-3} M_{\odot}\) based on the integrated area of the \(^{13}\text{CO}\) emission.

### 3.2.2. Compact \(^{12}\text{CO} 3–2\) and SO 5(6)–4(5) Emissions

The channel maps of \(^{13}\text{CO} 3–2\) and SO are shown in Figure 6 and the total intensity maps are presented in Figure 7. The \(^{12}\text{CO}\) 3–2 emission appears very compact and peaks at UY Aur A, either tracing the outflow(s), a disk, or just warmer molecular gas. In contrast to CO, SO is often found in outflow-driving and embedded sources (Guilloteau et al. 2013). The SO emission in UY Aur is seen in between the binary with two extensions: the main peak with extension toward the northwest and the other peak, hereafter, SO SE peak, at 1\(^{\prime}\)5 southeast to the main SO peak. In summary, SO and CO 3–2 emissions are more compact than \(^{13}\text{CO}\) and C\(^{18}\text{O}\) 2–1 emission and peak closer to the binary, likely tracing warmer gas than the \(^{13}\text{CO}\) and C\(^{18}\text{O}\).

#### 3.2.3. A Complex Velocity Pattern

The velocity structure of the UY Aur system is very complex, as shown in both the moment maps of \(^{13}\text{CO}\) and C\(^{18}\text{O}\) (Figure 8) and the position–velocity (PV) diagrams (Figure 9). The intensity-weighted velocity maps (moment 1, M1, Figure 8) of both \(^{13}\text{CO}\) and C\(^{18}\text{O}\) reveal that the main velocity gradient is along the major axis of the CB disk, suggesting that the velocity pattern is dominated by rotation at a large scale. Surprisingly, the velocity dispersion (moment 2, M2) is very large at the eastern edge of the CB ring and at UY Aur A. Although there is a general trend of redshifted gas in the northwest and blueshifted gas in the southeast, the system contains a significant amount of gas deviating from simple Keplerian motions.

The kinematics are better analyzed using the PV diagrams. In Figures 9(a) and (b), the PV diagrams are cut along the major axis of the disk (P.A. of 135\(^\circ\), shown as gray segments in Figure 8). In the blueshifted part, there are three gas components with velocity gradient: the outer most one is in Keplerian rotation and is probably tracing the extended CB disk, and the other two are close to the SO extension peaks. We note that there is about comparable emission strength between the extended CB disk and the ones associated with the SO extension peaks. Interestingly the PV diagrams along the binary (AB line) of 228\(^\circ\) (Figures 9(d) and (e)), which is almost perpendicular to the major axis of the disk, also show some high-velocity streamers in between UY Aur A and B. Such a velocity pattern is characteristic of infalling material as illustrated by the black dashed curve (\(V_{\text{infall}} \propto r^{-1}\)) in Panels (d) and (e) and likely traces gas coming from the CB ring that is accreted onto the individual disk components. There is also an outflow component detected at 1\(^{\prime}\)–2\(^{\prime}\) in the blueshifted part near UY Aur A (Figures 9(d) and (e)). The PV plots of \(^{12}\text{CO} 3–2\) along P.A. of 135\(^\circ\) and along P.A. 228\(^\circ\) are shown in Figures 9(c) and (f), respectively.
Contrary to CO 2–1 gas, the $^{12}$CO 3–2 only partly traces the CB ring (Figure 9(c), P.A. of 135°) and the infalling material (Figure 9(f), P.A. of 228°). It essentially traces warmer gas located near the stars and that is not seen in $^{13}$CO and C$^{18}$O 2–1 maps, which characterize colder material.

We further extract the spectra of $^{13}$CO 2–1, C$^{18}$O 2–1, and $^{12}$CO 3–2 at UY Aur A and B, at NIR 1, at the SO SE peak, and along the NW spiral (Figure 10). The locations where the spectra are taken along the NW spiral are marked in Figure 8. In general, the $^{13}$CO and C$^{18}$O show similar behavior, as
was seen in Sections 3.2.1 and 3.2.3. The $^{12}$CO 3–2 spectra have dramatically different shapes from $^{13}$CO and C$^{18}$O along NW Spi1 and NW Spi2, and at NIR 1. There is no good reference velocity to compare with the obtained spectra because the dominant force is highly dependent on the distances to UY Aur A or B. We note, however, that there are clearly high-velocity wings seen in $^{13}$CO in the redshifted end at the NW spiral, especially in Spi1 and Spi2.

### 4. DISCUSSION

#### 4.1. The Extended CB Ring

The CO emission is clearly dominated by the extended CB ring, but there is only a marginal detection of a dusty counterpart (Figure 1). This suggests that the amount of material in the CB ring is low, as estimated from both the $^{13}$CO 2–1 integrated emission and the 1.4 mm residual dust emission. The total mass (gas+dust) derived from the extended $^{13}$CO emission around UY Aur (hence including the streamers) is of the order of $M_{\text{gas}} \sim 2 \times 10^{-3} M_\odot$. The mass derived from the residual extended flux emission is $\sim 7 \times 10^{-4} M_\odot$. Hence the mass of the ring is of the order of $10^{-3} M_\odot$. This is two magnitudes lower than that of the ring orbiting GG Tau A Dutrey et al. (1994). Contrary to the GG Tau ring (Dutrey et al. 1994), where the inner edge of the CB ring is well defined, the inner edge of UY Aur is at $\gtrsim 540$ AU, in agreement with Duvert et al. (1998), and is not well defined. The contrast between the CB ring and the cavity between the CB ring and CS disks is moderate, as there is substantial gas emission in the cavity. There is about a comparable amount of emission inside the cavity as on the CB
ring (see Sections 3.2.1 and 3.2.3), making the determination of ring size, its sharpness, and its inclination more difficult.

Hioki et al. (2007) concluded that scattering anisotropy was insufficient to explain the strong contrast between the northwest and southeast parts of the CB ring, and that the far side of the ring was obscured by additional material. The anti-correlation between the NIR scattered light and the detected millimeter continuum emission provides direct support for this selective extinction hypothesis. The extended $^{13}$CO emission (see Figure 8) toward the east of the binary further supports this interpretation. We note that the kinematics is peculiar at NIR 1 peak. At the location of NIR 1, which is at the near side, there is an additional $^{13}$CO emission peak at redshifted velocity $V_{\text{LSR}}$ of 7.16 km s$^{-1}$ besides the blueshifted Keplerian rotation (Figure 10). This extended component has a size of $\sim$5$''$ with a projected velocity gradient of 1.4 km s$^{-1}$ (Figure 9). This cloud is likely gravitationally bound to the binary system since its velocity difference to the Keplerian speed at this location is only $\sim$0.6–1.4 km s$^{-1}$. As in the case of AB Auriga (Tang et al. 2012), such a feature may be a foreground CO cloud captured by the UY Aur system and falling down to the CB material, where the emission appears redshifted due to projection. This infalling gas at NIR 1 provides an additional explanation of the contrast of the CB ring at NIR.

4.2. Evidences for Streamers

Several extended features suggest the existence of accretion material between the CB disk and the CS disks. At NIR, a hint
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Figure 7. Integrated intensity images of $^{13}$CO 2–1 (a), C$^{18}$O 2–1 (b), SO (c), and 12CO 3–2 (d) overlaid on the NIR image (grayscale). In (a), the contour levels are 1, 2, 3, ..., 11 $\times$ 0.05 Jy beam$^{-1}$ km s$^{-1}$. The beam size is 1$''$27 $\times$ 0$''$92. In Panel (b), the contour levels are 1, 2, 3, ..., 12, 13 $\times$ 7 mJy beam$^{-1}$ km s$^{-1}$. The angular resolution is 1$''$94 $\times$ 1$''$74. In Panel (c), the contour levels are 6, 12, 18, 30, 36, 42 $\times$ 2.7 mJy beam$^{-1}$ km s$^{-1}$ (1$\sigma$). The beam size is 1$''$25 $\times$ 0$''$9. In Panel (d), the contour levels are 2, 4, 6, ..., 16, 18 $\times$ 1.1 Jy beam$^{-1}$ km s$^{-1}$ (1$\sigma$). The beam size is 1$''$12 $\times$ 0$''$98 with P.A. of 123$^{\circ}$.

Figure 8. Top panels: Moment 0, 1, and 2 maps of $^{13}$CO. The units of the wedge is in Jy beam$^{-1}$ km s$^{-1}$ for M0 and in km s$^{-1}$ for M1 and M2. The three small “+” in the panel M0 mark the positions NW Spi 1, 2, and 3 from west to east, where the spectra are taken. Gray segments mark the position–velocity cut along P.A. of 135$^{\circ}$ and along P.A. of 228$^{\circ}$ (shown in Figure 9). Bottom panels: the same as top panels but for C$^{18}$O. Note that the intensity of the C$^{18}$O M0 map is scaled by a factor of 6.4.

of possible spiral arm is identified in the J band by (Close et al. 1998). In the H band (Figure 1; Hioki et al. 2007), there are two features identified at the base of the inner edge of the ring, one of which corresponds to the possible spiral arm in the J band. An $H_2$ 2.12 $\mu$m arc of $\sim$0.5 in extension is also detected to the south of the CS disk of UY Aur A toward the east. Such an extension cannot be explained by the outflow since the bipolar outflow has a P.A. of 40$^{\circ}$ (Hirth et al. 1997). Although the excitation conditions are not yet determined, the detected $H_2$ arc may be due to accretion shocks onto the CS disk surrounding UY Aur A, as it is observed in the case of GG Tau Aa (Beck et al. 2012). In addition to the $H_2$ arc, diffuse $H_2$ emission is observed around both stars.

Our new detections of the SO, $^{12}$CO, $^{13}$CO, and C$^{18}$O gas surrounding the binary further support the existence of streamers. In the integrated intensity images, the northwest accreting
Figure 9. Velocity–position plots along P.A. of 135° centered on UY Aur A traced with $^{13}$CO 2–1 (a), C$^{18}$O 2–1 (b), and $^{12}$CO 3–2 (c). The solid curves mark the velocity of a Keplerian rotation with $i = 42°$ and an enclosed mass of 1.2 $M_\odot$. The dashed curves denote motions with conserved angular momentum ($V_{\text{rot}} \propto r^{-1}$). The cyan segments mark the SO peaks. The orange arrow marks the NIR 1 location. The resolutions are marked in the top right corner. The plots along the UY Aur A and B (P.A. = 228°) traced with $^{13}$CO (d), with C$^{18}$O (e), and with $^{12}$CO 3–2 (f). The blue segments mark the location of UY Aur A and B. The black dashed curves denote the free-fall motion with $i = 42°$ and an enclosed mass of 1.2 $M_\odot$. In all panels, $V_{\text{sys}}$ is 6.2 km s$^{-1}$ (shown in dotted line). (A color version of this figure is available in the online journal.)
Figure 10. Spectra of $^{13}$CO 2–1, C$^{18}$O 2–1, and $^{12}$CO 3–2 toward UY Aur A, UY Aur B, the NIR 1 peak, the SO SE peak, and three spots of the NW Spiral marked as "+" in Figure 8. The vertical dotted line marks the $V_{\text{sys}}$ of 6.2 km s$^{-1}$. The dashed lines mark the 3$\sigma$ level.

material is clearly resolved in $^{13}$CO and C$^{18}$O (Figure 7). Both in $^{13}$CO and C$^{18}$O 2–1, the cavity does not appear free of gas. At the inner edge of the CB disk, the emission is asymmetric (Figure 8). Moreover, at the location of the SO extension peaks, there is some CO gas moving with a velocity gradient of 1–2 km s$^{-1}$ traced in $^{13}$CO and C$^{18}$O (see the PV plots in Figures 9(a) and 9(b); spectra in Figure 10). In the two lines, confusion along the line of sight due to the NIR 1 component also precludes a good identification of the SE spiral. Contrary to $^{13}$CO and C$^{18}$O 2–1, the $^{12}$CO 3–2 emission is more compact and confined within the gap of the CB disk (Figure 7). Such a difference, likely due to a temperature effect in the ring, also suggests that there are gradients of excitation conditions inside the ring.

The SO emission peaks around the binary with east and west extensions along the axis perpendicular to the binary. We also note that the detected SO extensions are closer to the primary UY Aur A instead of the secondary. SO is known to be a shock tracer. Although we cannot reject that some SO emission may come from the outflow, the extended SO emission may come from a shock point at the interface between the inner edge of the CB disk and the inward accreting material, as predicted by, e.g., Hanawa et al. (2010), where the gas inflows from the CB disk through L2. Based on simulations and depending on the binary orbit, its eccentricity, and the mass ratio, this gas is further captured mainly by the primary (e.g., Ochi et al. 2005) or by the secondary (Bate & Bonnell 1997). Future ALMA observations with multi-J transitions of the CO isotopologues and the SO lines at higher angular resolution would provide a clearer picture of the streamers and their links with both the CB and CS material.

4.3. Material Surrounding the Individual Stars

When UY Aur was identified as a binary by Joy & van Biesbroeck (1944), the two stars were optically visible. UY Aur B has been once classified as an infrared companion (Tessier et al. 1994). Its magnitude changed on a long period by several magnitudes since 1944 (Berdnikov et al. 2010). Interestingly, several observational evidences also indicate that the accretion process around the binary varies significantly within one decade. Skemer et al. (2010) report that both UY Aur A and B varied at mid-infrared between 1999 to 2009, and that the dust grains surrounding UY Aur A are extremely pristine and ISM-like based on the silicate spectrum. Skemer et al. (2010) interpret such dust properties as an indirect proof of replenishment by a less evolved CB dust through spiral streamers. Similarly, around UY Aur B, Herbst et al. (1995) argued that such changes in magnitude cannot be caused by spherically distributed material, but may be caused by inhomogeneous accretion from a flattened structure. Moreover, Herbst et al. (1995) reported the existence of an outflow while Pyo et al. (2014) revealed that both UY Aur A and B drive optical jets along an axis at P.A. of $\sim$40$^\circ$, in agreement with the outflow axis. Finding a scenario that reconciles all particularities of the binary system seems to be difficult. In this section, we discuss some facts, their implications on the system, and some possible ways to reconcile all the observed properties of the source.

4.3.1. UY Aur A

We first note that the low spectral index of the millimeter emission, $\sim$1.6, can only be explained either by optically thick dust emission at low temperatures (10 K with an upper limit of 20 K given the uncertainties), or by a superposition of warmer dust and free–free emission at longer wavelengths. As the emission originates from distances smaller than 20 AU from the star, the required temperature appears unrealistically low. Free–free emission from an ionized jet is expected to have a spectral index of 0.6, while dust at 25 K would have an index of 1.9. The kink in the spectral energy distribution indicated by the 0.85 mm flux measurement made with ALMA by Akeson & Jensen (2014) further suggests a contamination by free–free emission, which becomes negligible only above 300 GHz. In this scenario, we may attribute 18 mJy to the dust emission at 220 GHz, which requires a minimum radius of 12 AU using the typical compact disk properties as described in Piétu et al. (2014).

Berdnikov et al. (2010) discussed the $B$ and $V$ band light curves on a timescale of 100 yr as the infrared variations of UY Aur A and B. They identified a periodicity of 16.3 yr in
the variations of UY Auriga A, which they interpreted as a modulation of the accretion rate induced by a companion located at 6 AU from the star. Our detection of millimeter continuum emission restricts the possibilities for such a companion. The invoked companion should not be too massive in order to avoid opening a substantial gap in this disk. Placing a firm upper limit on the body mass is premature, however, as the width and depth of tidal gaps also depend on the disk viscosity (Crida et al. 2006).

4.3.2. UY Aur B as a Binary?

Besides the low spectral index, the millimeter emission around UY Aur B has two further unusual characteristics: (1) it is elongated along the overall UY Aur system rotation axis and (2) it is offset by 0′07 (deprojected, 10 AU) from the optical position of UY Aur B along the projected system axis. A simple way to explain the elongation and shift is to assume this is due to superposition of two equally bright and unresolved features separated by 0′14, or 20 AU. This suggests that B could itself be a close binary, each with a dusty disk emitting at millimeter wavelengths.

The long period of decreased optical magnitude for UY Aur B, of duration $D_d \approx 40$ yr between approximately 1960 and 2000, would be due to extinction of UY Aur Ba by the disk of the less massive companion UY Aur Bb. The orbiting period $P_{Bb}$ of the Ba/Bb binary can be estimated from $(2\pi/P_{Bb}) = (2\theta/D_d)$, where $\theta = \text{asin}(R_{b}/a)$ and $R_{b}/a$ is the ratio of Bb disk radius to orbit semi-major axis $a$, $R_{b}/a \leq 1/3$ because of tidal truncation. Using $D_d = 40$ yr yields a period $P_{Bb}$ of 360 yr. Assuming a total mass $M_{(a+b)}$ of 0.4 $M_{\odot}$ (Hartigan & Kenyon 2003), the separation between Ba and Bb will be 37 AU if the orbit is circular. This orbital separation is larger than the deprojected separation of 20 AU from the assumption of equal masses earlier in the section. Such a discrepancy can be explained if the inclination angle is larger than 42°. If the inclination is up to $\sim 66°$, these two estimated values will be consistent. More reasons in favor of a larger inclination of the UY Aur system are discussed in Section 4.4.

The flux from each component, around 3.4 mJy each at 1.4 mm, could be explained by very compact, $\approx 5$ AU, essentially optically thick disks. A spectral index of 1.6–1.8 requires $T < 20$ K, which however looks very small for a radius of 5 AU, even though the two stars would be rather low-mass objects. An alternative is optically thin disks of radii about 20 AU, so that their mean temperature does not exceed 20 K and with grains large enough to have an emissivity exponent $\beta = 0$. This interpretation is thus marginally plausible in terms of sizes and timescales. Moreover, in this hypothesis, UY Aur A, Ba, and Bb are nearly aligned with the projection of the system axis, so an obscuration is only possible if the orbit of Ba/Bb is perpendicular to the rest of the system. Yet the optical jet emanating from B recently detected by Pyo et al. (2014) is aligned on the overall system axis, which suggests at least one of the Ba/Bb disk is coplanar with the overall UY Aur system.

4.3.3. Jets

We already argued in Section 4.3.1 that free–free emission from an optical jet provides an appropriate explanation to the low spectral index of the millimeter continuum emission from UY Aur A. UY Aur B displays a similar spectral index, although the error bar is larger. It is thus also plausible that the emission from UY Aur B is also the superposition of dust emission from a CS disk with free–free emission from an ionized jet. This jet is aligned along the system axis (Pyo et al. 2014). This can offer a simple interpretation for both the source elongation and the position offset, provided only the southeast jet emits in free–free. Indeed, Pyo et al. (2014) have already shown a strong asymmetry in the [Fe II] emission, which is only detected toward the northwest of UY Aur B, while it is detected on both sides of UY Aur A. In addition, an elongated jet offers more flexibility to explain the relatively large size of $\sim 0′.45$ compared to the simple superposition of two point sources for which it is difficult to obtain a size above $\sim 0′.3$ for a centroid shift of $0′.07$.

4.4. A More Inclined System?

While the “disks in a binary” scenario (Section 4.3.2) has difficulty explaining the jet orientation, the “jet only” scenario (Section 4.3.3) unfortunately offers no explanation for the long-duration fading of UY Aur B. We explore here a possibility to combine both aspects and provide a consistent view of all phenomena affecting UY Aur.

The key difficulty of the binary scenario for UY Aur B is the near alignment of A, Ba, and Bb in combination with the rather low inclination of the system. On one hand, given the positional uncertainty, the alignment constraints between the A–B and Ba–Bb lines can be relaxed to about 45°. On the other hand, the inclination angle of 42°, which was first given by Duvert et al. (1998), is rather uncertain. Hioki et al. (2007) quote 42° $\pm 3°$ from a fit of an ellipse to the brightest spots in the NIR scattered light image. However, the fit of the fitted ellipse does not go through the center of mass of the system, which must be in between A and B, about twice as close to A than B. Hioki et al. (2007) also pointed out that the phase function of scattering cannot be anisotropic enough to explain the northwest/southeast contrast with this inclination and some asymmetric extinction is needed to explain northern–southern difference.

The dynamical mass of 1.2 $M_{\odot}$ estimated by Duvert et al. (1998) for $i = 42°$ becomes 0.7 $M_{\odot}$ for $i = 60°$. This is small compared to the latest estimate of the stellar masses from the spectral type by Hartigan & Kenyon (2003), being 0.60 and 0.34 $M_{\odot}$ for A and B, respectively. We note that the identification of a Keplerian pattern in the emission from the CO isotopologues is difficult, so the dynamical mass is rather uncertain. However, Hioki et al. (2007) derived a total binary mass to be 1.73 $\pm 0.29$ $M_{\odot}$ using the 42° inclination. This is already above the 0.94 $M_{\odot}$ value, and should scale as 1/$\cos^3(i)$ since the A–B separation is nearly aligned with the projected minor axis. A higher inclination can thus be accommodated only if the orbit is elliptical and the stars near apastron.

Hence, a higher inclination of $\sim 50°$–60° for the overall system, combined with inclination differences up to 10° between the various disks and orbital planes, may be sufficient to allow obscuration of Ba by Bb. The stability of such a system has also to be investigated. Furthermore, while Ba is exiting the occultation by Bb, its southern–eastern jet may still be behind the Bb disk, offering a direct explanation for the lack of detection of [Fe II] in this direction by Pyo et al. (2014).

We note that these differences in inclination are reminiscent of the situation in AB Aur, where the presence of surrounding material leads to an apparent inclination quite different from that of the inner disk (Tang et al. 2012; Piétu et al. 2005).
The surroundings of UY Aur are complex and may still influence the evolution of the system.

Figure 11 is a scheme which summarizes most of the observed properties of the UY Aur system. The accreting material between the CB and CS disks is better demonstrated in the PV diagrams (Figure 9). The proposed scenario (binary) can be tested by searching for a binary companion (Ba) at 40 AU from B (Ba) using NIR interferometric techniques. Similar results have recently shown by Di Folco et al. (2014), who found that GG Tau A is in fact a triple system, the second binary located around GG Tau Ab having a separation of 5 AU. If the discussed scenario is valid, the UY Aur should be a triple system.

5. CONCLUSION AND SUMMARY

Using the IRAM array and the SMA, we have studied the properties of the molecular gas and dust orbiting around the binary system UY Aur. We observed the CB ring, the streamers, and the CS disks with the molecular lines $^{13}$CO and C$^{18}$O 2–1, SO 5(6)–4(5) and $^{12}$CO 3–2. The thermal dust emission at 3.4 mm, 1.4 mm, and 0.85 mm is also studied. The results are summarized as follows.

1. The gaseous component of the CB ring is easily detected in $^{13}$CO and C$^{18}$O but its dust counterpart is very marginally seen, even at 1 mm continuum. We derive a total mass of $M_{\text{gas}} \sim 2 \times 10^{-3} M_\odot$. A gaseous component associated with NIR 1 in the southeast of the CB ring is detected, which is probably infalling toward the CB ring.
2. The inflowing cold gas emanating from the CB ring toward the CS disks is clearly detected in $^{13}$CO and C$^{18}$O following spiral patterns in the northeast and in the southwest of the CB disk, revealing that accretion proceeds through streamers. Combinations of a Keplerian CB ring, inflowing gas, and also probably outflows make a complex velocity pattern and preclude any accurate determination of the mass of the central binary.
3. The detections of $^{12}$CO 3–2 and SO in the vicinity of the CS disks suggest gradients of physical conditions in the streamers and likely accretion shocks, at least near the CS disk of UY Aur A. This scenario is also supported by the detection of bright $H_2$ patches near UY Aur A.
4. At least two CS disks are well detected in the millimeter continuum emission. The spectral indices $\alpha$ are 1.62 ± 0.1 and 1.62 ± 0.3 for UY Aur A and B, respectively. Such low spectral indices are more likely due to the combination of dust emission and free–free continuum emission.
5. To reconcile the optical variability of UY Aur A with the existing bright millimeter dust disk, we propose the object orbiting UY Aur A in close orbit (6 AU) is a planet whose mass is not enough to open a large gap in the disk around UY Aur A.
6. To reconcile the optical and infrared variability of UY Aur B, we propose this object is a binary of separation of the order of 40 AU (to take into account the long duration of the obscuration ~40 yr). Such a binary scenario is compatible with the presence of a jet but requires a higher inclination for the overall system (55–60° instead of 42°), combined with inclination differences as high as 10° between the various disks and orbital planes.

In conclusion, these observations reveal that if accretion proceeds through streamers in binary systems, the streamer configurations, their strengths, and their dynamics are difficult to
constrain. They have to be investigated in detail to confront models to the observations, requiring the sensitivity and resolving power of ALMA complemented by optical (interferometric) observations.

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Facilities: IRAM:Interferometer, SMA

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