ABSTRACT: Extracellular histone H4 is an attractive drug target owing to its roles in organ failure in sepsis and other diseases. To identify inhibitors using in silico methods, information on histone H4 structural dynamics and three-dimensional (3D) structural coordinates is required. Here, DNA-free histone H4 type 1 (H4.1) was characterized by utilizing tandem nonlinear and linear ion mobility spectrometry (FAIMS-TIMS) coupled to mass spectrometry (MS) complemented with molecular dynamics (MD) simulations. The gas-phase structures of H4.1 are dependent on the starting solution conditions, evidenced by differences in charge state distributions, mobility distributions, and collision-induced unfolding (CIU) pathways. The experimental results show that H4.1 adopts diverse conformational types from compact (C) to partially folded (P) and subsequently elongated (E) structures. Molecular dynamics simulations provided candidate structures for the histone H4.1 monomer in solution and for the gas-phase structures observed using FAIMS-IMS-TOF MS as a function of the charge state and mobility distribution. A combination of the FAIMS-TIMS experimental results with theoretical dipole calculations reveals the important role of charge distribution in the dipole alignment of H4.1 elongated structures at high electric fields. A comparison of the secondary and primary structures of DNA-free H2A.1 and H4.1 is made based on the experimental IMS-MS and MD findings.

INTRODUCTION

Histones are intrinsically disordered proteins\(^1\) that package eucharyotic DNA into nucleosomes within cell nuclei.\(^2,3\) A nucleosome consists of 147 base pairs of DNA wrapped around an octameric core histone comprising two pairs of each heterodimer of histones H3 and H4 and histones H2A and H2B.\(^2,3\) A nucleosome is connected to its adjacent nucleosome by a DNA segment associated with a linker histone H1 or H5 to form chromatin fiber.\(^4\) A core histone contains a central helix fold and an intrinsically disordered and charged N-terminal tail, which is protruded from the nucleosome and is highly subjected to post-translational modifications (PTMs) such as acetylation, methylation, phosphorylation, and glycosylation.\(^2,5,6\) The functions of histone PTMs are of particular interest and have been investigated extensively because combinatorial PTMs create “histone codes”\(^5\) that regulate DNA transcriptional activity,\(^5,7\) DNA replication,\(^5\) DNA repair, and recombination.\(^8\) Moreover, alterations in patterns of histone PTMs have been related to deregulation of the expression of many genes having crucial functions in cancer development and progression.\(^6,9\) During cell metabolism, newly synthesized DNA-free histones in the cytoplasm are transported into the nucleus by chaperones for replacement of old/damaged histones.\(^10\) Histones can be released into the extracellular space after significant cellular death such as sepsis, trauma, ischemia-reperfusion injury, and autoimmune disease.\(^11\) Once entered the extracellular space, histones function as damage-associated molecular pattern proteins, triggering the immune system and causing cytotoxicity.\(^11\) Extracellular histones can be found in a form of nucleosome-containing structures, neutrophil extracellular traps (NETs, comprising ~70% histones), or DNA-free histones.\(^12\) Degradation of NETs and nucleosomes in the bloodstream releases DNA-free histones, which are cytotoxic due to their highly cationic nature.\(^12\) Recent studies have shown that DNA-free histones play important roles in organ failure and death in sepsis and other diseases.\(^12\) In addition, a study in COVID-19 patients showed a correlation of the serum-citrullinated histone H3 level with acute respiratory distress syndrome.\(^12,15\) Moreover, it has been reported that histone H4 functions as a mediator causing arterial tissue damage and inflammation.\(^16\) Furthermore, histone H4 also induces platelet aggregation in thrombocytopenia.\(^17\) These studies imply a new trending strategy in the discovery of novel drugs that effectively inhibit the release of histones and/or neutralize circulating histones.

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Recently, small polyanions and peptidic inhibitors have been developed, targeting electrostatic interactions with histone H4 and H2A monomers to neutralize their pathological effects. The successes of these studies represent a promise for the development of inhibitors targeting intrinsically disordered histone monomers. With the advent of advanced molecular modeling software, in silico drug discovery approaches have been successfully applied for identification of potential lead inhibitors before extensive experimental assays. To identify inhibitors using in silico methods, information on the protein intrinsic dynamics and their three-dimensional (3D) structures is required. However, the structural analysis and 3D structures of the H4.1 monomer are currently absent. Furthermore, for efficient identification and design of inhibitors for histone monomers using in silico approaches, a proper understanding of biomolecular recognition mechanisms is needed. For instance, protein flexibility such as conformational changes prior and upon association with inhibitors is diverse and not obvious when only a few or no protein structures are available. In addition, using one protein conformation for structure-based screening restricts the chemical space of potential inhibitors for a specific drug target. Therefore, an ensemble of protein structures needs to be exploited to improve the accuracy.

In this paper, we applied the recently developed methods and strategies for successful structural analysis of an intrinsically disordered histone H2A.12 to characterize the gas-phase structures of human histone H4.1. Briefly, in nonlinear ion mobility spectroscopy (i.e., field asymmetric IMS, FAIMS), ion mobilities are separated based on different K values (ΔK) between the low and high electric fields (E), whereas in linear IMS, particularly trapped IMS (TIMS), ion mobilities are measured based on the rotationally averaged collision cross section (CCS). The IMS-mass spectrometry (IMS-MS) method coupled to nanoelectrospray ionization (nESI-IMS-MS) has demonstrated significant advances for investigation of protein conformational states owing to its capability to retain solution structures upon transfer to the gas phase, also known as the “solution memory effect”. In addition, ions of the protein at high charge states or upon activation are associated with partially unfolded and/or extended structures, analogous to the unfolding processes induced by protonation in solution. Furthermore, the integration of experimental ion mobilities (K) with theoretical calculations for candidate structures generated from molecular dynamics (MD) simulations has established IMS as a powerful structural probe.

This report illustrates novelty and originality, as it provides gas-phase candidate structures from native-like to extended conformations as well as unfolding pathways for the histone H4.1 monomer. The findings of this study will provide the starting structure of the histone H4.1 monomer and information on the protein flexibility for the efficient discovery of specific inhibitors for the H4.1 monomer.

## RESULTS AND DISCUSSION

**nESI-TIMS-FAIMS-TOF MS.** The MS spectra of H4.1 using native (65 mM NH₄Ac and 10 mM NH₄Ac) and denatured starting solutions (organic content) acquired using the nESI-TIMS-TOF MS configuration are shown in Figure 1a–c, respectively. The MS spectrum of native H4.1 exhibited a wide charge state distribution (CSD) spanning from charge state (z)

![Figure 1. TIMS-MS mass spectra (a–c) and CCS profiles (d–f) of native H4.1 in 65 mM NH₄Ac, 10 mM NH₄Ac, and denatured H4.1, respectively. Data were acquired using "soft" TIMS ion injection settings with Srf,TIMS = 1.6 V/ms. The CCS profile of [M + 5H]⁺ is not shown due to a low signal. Note that CCS values increase with increasing charge states. The red lines mark a trend for C, P, and E conformations of H4.1.](https://doi.org/10.1021/acsomega.1c03744)
injection ($\Delta V_{\text{TIMS}} = 20$ V). Inspection of these TIMS-MS CCS profiles revealed an increase in the CCS with the increase of charge states with multiple CCS bands observed per charge state in most of the intermediates and higher charge states. For example, the CCS profile of native H4.1 in 65 mM NH$_4$Ac showed that $z = 6^+$ and $7^+$ exhibited a bell shape centered at $\sim$1470 and 1583 Å$^2$, respectively. These bands are likely characteristics of compact-like (C) conformations. The CCS profile of $z = 8^+$ displayed a major band centered at $\sim$1679 Å$^2$ and minor bands at 1792, 1935, and 2395 Å$^2$, indicating the coexistence of C-like structures with a slightly less compact conformation. The CCS profiles of $z = 9^+$ to $13^+$ display many CCS bands, consistent with coexisting ions populating C-like and partially unfolded (P-like) structures. The CCS profiles of $z = 14^+$ to $17^+$ display a narrow CCS band with maxima at 3220, 3272, 3360, and 3470 Å$^2$, respectively. These characteristics imply that native H4.1 can adopt elongated (E-like) structures at higher charge states. When compared with the TIMS-MS CCS profiles of native H4.1 in 10 mM NH$_4$Ac, no differences were observed; however, considering better ion signal intensity observed in 10 mM NH$_4$Ac, this condition was used for native nESI-FAIMS-TIMS experiments. The TIMS-MS CCS profiles of denatured H4.1 (Figure 1f) showed some noticeable differences compared to those of native H4.1. For example, more CCS bands at $\sim$1682, 1720, 1815, 1952, 2210, 2316, and 2431 Å$^2$ were observed for $z = 8^+$ of the denatured form (Table S1). Additionally, these CCS bands are more pronounced for denatured H4.1 compared to their native counterparts. These results further suggest “solution memory effects” on gas-phase structures of H4.1, evidenced not only by changes in CSD but also by mobility profiles.

A comparison of CCS profiles for native (10 mM NH$_4$Ac) and denatured H4.1 acquired using nESI-TIMS-TOF MS (blue traces) and nESI-FAIMS-TIMS-TOF MS (red traces) configurations at soft ion injection ($\Delta V_{\text{TIMS}} = 20$ V) is shown in Figures S1b and S2b. The FAIMS-TIMS-MS CCS profiles of the lowest ($z = 6^+$ and $7^+$) and highest ($z = 14^+$ to $17^+$) charge states were close to those obtained using the TIMS-MS
configuration. The FAIMS-TIMS-MS CCS profiles of charge states 9+ to 13+ follow the established patterns of TIMS-MS CCS profiles with the coexistence of C- and P-like structures. A closer inspection revealed differences in the CCS profiles of data acquired using TIMS-MS and FAIMS-TIMS-MS configurations. That is, the addition of the FAIMS stage significantly affected the CCS profiles of the P-like structures (z = 8+ to 11+), consistent with our previous report. For instance, in the case of z = 9+, the major CCS band was shifted to a higher CCS value in FAIMS-TIMS-MS (∼2020 Å²) compared to TIMS-MS (∼1880 Å²). A similar trend was observed for z = 8+, 10+, and 11+. These results implicate a subtle unfolding induced by ion heating in the FAIMS stage, consistent with previous reports.  

The compensation field profiles (E_c) of FAIMS-TOF MS (back traces) and FAIMS-TIMS-TOF MS (red traces) for H4.1 in native (10 mM NH4Ac) and denatured starting solutions are shown in Figures S1c and S2c, respectively. No major difference in E_c profiles was observed between the data of native and denatured conditions, as well as with the inclusion of the TIMS separation. The total ion chromatogram (TIC) comprises ion signals ranging from −5 V/cm ≤ E_c ≤ +55 V/cm, suggesting the existence of similar structures with varying dipole alignments. Overall, the E_c profiles of H4.1 exhibited higher E_c values for C-like structures (z = 6+ and 7+), while a decreasing trend toward −E_c represents P-like and/or E-like structures. We noticed that a bimodal distribution with the lowest E_c values was observed for E-like structures, consistent with our previous report for histone H2A type 1 (H2A.1).  

To correlate features of CCS with E_c, the two-dimensional (2D) palettes of H4.1 were generated. Features in the 2D palettes of native and denatured H4.1 showed two features A1 (CCS of ∼2020 Å² and E_c of ∼40 V/cm), B1 (CCS of ∼2200 Å² and E_c of ∼33 V/cm), A2 (CCS of ∼2470 Å² and E_c of ∼40 V/cm), B2 (CCS of ∼2500 Å² and E_c of ∼20 V/cm), and B2 (CCS of ∼2470 Å² and E_c of ∼20 V/cm). These features are characteristic of C-like (feature A1), P-like (features B1 and A2), and E-like (features C1 and C2) structures. The 2D palettes of z = 10+ native and denatured H4.1 showed two features A1 (CCS of ∼2280 Å² and E_c of ∼38 V/cm) and A2 (CCS of ∼2500 Å² and E_c of ∼38 V/cm) at a higher E_c range, corresponding to P-like structures. In contrast, feature B1 (CCS of ∼2500 Å² and E_c of ∼15 V/cm) was observed at a lower E_c range, consistent with E-like structures due to more dipole alignments. The 2D palettes of z = 11+ of native and denatured H4.1 follow the established patterns of z = 10+. We noticed that features A2 and B1 of z = 10+ and 11+ were more abundant for the native form. For z = 12+ and 13+, native and denatured H4.1, features A1, A2, and A3 have similar E_c values but significantly different CCS values, corresponding to P- and E-like structures. In addition, feature B1 was observed at a lower E_c range, corresponding to E-like structures with more dipole alignments. For z ≥ 14+, the 2D palettes of native and denatured H4.1 exhibited features A1 and A2 with slightly different CCS values in a higher E_c range, corresponding to the presence of multiple E-like structures. Moreover, feature B1 of these charge states exhibits higher relative abundance than that of the lower charge states, consistent with more dipole alignments.  

Collision-Induced Unfolding of H4.1. Our TIMS-MS and FAIMS-TIMS-MS results revealed that H4.1 adopts heterogeneous structures. To further characterize the gas-phase structural features of H4.1, we utilized collision-induced unfolding (CIU) to probe the conformational change as a function of activation voltage. The CIU fingerprints for charge state 7+ of native and denatured H4.1 are shown in Figure 3. The CIU fingerprints of all charge states are displayed in Figure S3.  

The inspection of H4.1 CIU fingerprints revealed a significant impact of the native (10 mM NH4Ac) and denatured starting solutions on the gas-phase H4.1 structures. The major differences were observed for the lowest charge state. For example, the CIU fingerprint of z = 7+ (Figure 3) for both native and denatured H4.1 exhibited ions populating C- and P-like structures; however, their CIU pathways were dependent on the starting solution conditions. A closer inspection of the fingerprint for native H4.1 (Figure 3a) revealed that at the lowest collision voltage (∼5−30 V), the CCS was centered at ∼1580 Å² (state 0), corresponding to C-like structures. As the collision voltage was increased (∼25−130 V), the CCS dramatically increased to ∼1870 Å².
These results are consistent with collision-induced unfolding.37 Interestingly, at higher collision voltages (∼125 to 220 V), the CCS decreased to ∼1750 Å² (conformation 2 of state 1), implying the refolding of H4.1 toward a more compact conformation. Typically, gas-phase activation results in unfolding;37 however, compaction of protein structures upon activation has been reported.21,38,39 This is particularly relevant since H4.1 has an intrinsically disordered N-terminal tail, which may or may not be exposed during the gas-phase unfolding transient. We noticed that a low abundance of conformation 1 (CCS of ∼1870 Å²) remained at the highest collision voltages (∼160 to 220 V), suggesting the coexistence of conformations 1 and 2.

A closer inspection of the CIU fingerprint for denatured H4.1 (Figure 3b) showed that at the lowest collision voltage (∼5 to 25 V), the CCS value was observed at ∼1580 Å² (state 0 at low CIU), corresponding to C-like structures. When the collision voltage was increased from 10 to 120 V, the CCS increased to ∼1650 Å² (conformation 1 of state 1), consistent with P-like structures. As the collision voltage was further increased from 25 to 220 V, the CCS increased to ∼1690 Å² (conformation 2 of state 1). In addition, conformation 3 with higher CCS (∼1870 Å²) coexisted in the same voltage range. When the collision voltage was increased from 50 to 125 V, the CCS further increased to ∼1910 Å² (conformation 4 of state 1). Interestingly, by further increasing the collision voltage from 125 to 200 V, the CCS decreased to ∼1790 Å² (conformation 5 of state 1), consistent with more compact structures. At the highest collision voltage (∼150 to 220 V), the CCS further decreased to ∼1580 Å² (state 0 at high CIU), corresponding to C-like structures. These results suggest that while the denatured starting solution maintained H4.1 compact structures, it also induced intramolecular arrangements to differentiate the CIU patterns of native and denatured H4.1. In the case of z = 8+, the CIU fingerprints of native and denatured showed mainly three CIU states: 0 for C-like, 1 for P-like, and 2 for E-like structures. In addition, no difference was observed for the CIU fingerprints for native and denatured H4.1 at higher charge states (z = 9+ to 17+).

CMD Simulations of Solution-Phase H4.1. Figure 4a shows the amino acid sequence, α-helices, 310-helices, and the N- and C-terminal regions of H4.1. The radius of gyration (R<sub>gyr</sub>) profile along the simulation time was utilized for the analysis of H4.1 structural rearrangement in solution. For example, the snapshots (A–I) of H4.1 conformations are shown in Figure 4b. The initial DNA-bound X-ray structure of H4.1 (A) comprises three α-helices (α1 to α3), one 310-helix, and an intrinsically disordered N-terminal tail. A close inspection of conformation A revealed that H4.1 structural elements are stabilized in the nucleosome by hydrogen bonding to DNA and by interactions with α-helices from histones H3.3 In addition, the N-terminal tail of the DNA-bound H4.1 structure was highly extended, leading to an extended conformation with a calculated R<sub>gyr</sub> of 19.4 Å. At 16 ns, the N-terminal tail was partially folded via formation of the R23–E63 salt bridge (conformation B). At 115 ns, the N-terminal vicinity was collapsed due to melting of the 310-helix

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**Figure 4.** Molecular modeling. (a) Amino acid sequence of H4.1. Residues corresponding to α- and 310-helices and the N- and C-terminal regions are marked. (b) Radius of gyration as a function of CMD simulation time. The snapshots of the solution phase of H4.1 are labeled from A to I. Residues corresponding to salt bridges are indicated. (c) Comparison of the gas-phase unfolding pathways between H4.1 and H2A.1.
and the rearrangement of α3 on the C-terminus, leading to formation of conformation C ($R_{gyr} = 20.1$ Å). At 140 ns, the $3_{10}$-helix was gradually refolded, and the N- and C-terminal tails were folded, resulting in conformation D ($R_{gyr} = 18.8$ Å), stabilized by salt bridges K$^{+}$/E$^{-}$, K$^{+}$/E$^{-}$, and D$^{-}$/R$^{+}$. At 215 ns, the N-terminal tail was partially unfolded, while the C-tail was folded, leading to conformation E ($R_{gyr} = 19.1$ Å). Up to 295 ns, both the N- and C-terminal tails were folded, leading to conformation F with a $R_{gyr}$ of 15.5 Å, stabilized by salt bridges K$^{+}$/D$^{-}$ and R$^{+}$/D$^{-}$. At 320 ns, the N-terminal tail was unfolded, leading to partially unfolded conformation G ($R_{gyr} = 17.6$ Å) stabilized by salt bridges R$^{+}$/E$^{-}$, R$^{+}$/E$^{-}$, K$^{+}$/E$^{-}$, and R$^{+}$/D$^{-}$. At 321 ns, both the N- and C-terminal tails were folded, leading to compact conformation H with the smallest $R_{gyr}$ of 14.8 Å, stabilized by salt bridges R$^{+}$/E$^{-}$, K$^{+}$/E$^{-}$, R$^{+}$/D$^{-}$, and K$^{+}$/E$^{-}$. Interestingly, the $\alpha$2-helix was slightly bent near the N-terminal vicinity (H, Figure 4b). Further simulations (e.g., $t = 495$ ns) lead to conformation I with a larger $R_{gyr}$ of 17.2 Å due to unfolded N-terminal tail. The effects of water box thickness and nonbonded cutoff are shown in Figure S5 and described in the Supporting Information.

**Comparison of Gas-Phase Structural Features between H2A.1 and H4.1.** Figure 4c shows the representative gas-phase candidate structures for the main CCS features of H4.1 observed from TIMS-MS experiments (all candidate structures for 6+ to 18+ are shown in Figure S4 and summarized in Table S1). For the most compact structures (6+ to 9+), two trends were observed: (i) a gas-phase compact structure where the $\alpha$2-helix was greatly bent and partially melted and (ii) a gas-phase compact structure that preserves the $\alpha$2-helix structure from that of the native solution (conformation H). A closer inspection revealed that the major structural changes in H4.1 were associated with the unfolding of N- and C-terminal tails, disruption of inter-α-helical interactions, leading to exposed charge residues, increasing the distance of N- and C-terminal domains, and extending the N- and C-terminal tails in higher charge states. For example, the candidate for $z = 8+$ exhibited a C-like structure with a calculated CCS of 1877 Å$^2$. A comparison of the unfolding pathways between H4.1 and H2A.1$^{11}$ illustrated a potential unfolding pathway for H4.1 as a function of charge state and CCS ranges. The candidate structures of both H4.1 and H2A.1 in the C-like conformation have the N- and C-terminal tails tightly folded toward the $\alpha$-helical core. The candidate for $z = 9+$ of H4.1 displayed a subtle rearrangement in both N- and C-terminal domains, leading to a refolding of the $3_{10}$-helix in the N-terminal vicinity and a slightly unfolding of the C-terminal tail, consistent with the structural rearrangement of the candidate for $z = 9+$ of H2A.1. Overall, these results demonstrate that the C-like structure of H4.1 exhibits a similar folding motif as observed for H2A.1. That is, both N- and C-terminal domains were folded toward the $\alpha$-helical core. The candidate structure of $z = 10+$ displays a major structural rearrangement associated with melting of $\alpha$-helices, leading to the destabilization of the secondary structure and exposed basic and acidic residues K$^{+}$, R$^{+}$, K$^{+}$, R$^{+}$, K$^{+}$, R$^{+}$, K$^{+}$, R$^{+}$, K$^{+}$, and E$^{-}$, consistent with the unfolding of H2A.1. Inspection of candidate structures of $z = 11+$ to 18+ (Figures 4c and S4) revealed that H4.1 adopted E-like structures with the high content of the random coil, in agreement with our experimental results. For example, most of the helical interactions in the candidate structure of $z = 11+$ were loose, followed by the distance between the N- and C-terminal domains. Moreover, the candidate structures of $z = 12+$ and 13+ exhibit a further expansion between the N- and C-terminal domains, leading to a highly extended conformation. Furthermore, the structures of $z = 13+$ and 14+ display a rearrangement of the N- and C-terminal domains, promoting an E-like conformation. For $z \geq 15+$, the candidates exhibit highly random coil structures with an extended C-terminus.

**Influence of the Charge Distribution on the Dipole Alignment of H4.1.** The 2D FAIMS-TIMS palettes present a correlation between dipole alignment, CCS, and the charge state distribution. The effect of charge on the dipole alignment is illustrated by charge state 17+ of denatured H4.1. As shown in Figure S$a$, the candidate structures representative for features A1 (CCS of $\sim$3402 Å$^2$) and A2/B1 (CCS of $\sim$3442 Å$^2$) adopt an E-like conformation with a relatively narrow CCS distribution. In addition, these candidates have high structural similarity, evidenced by a small backbone RMSD of $\sim$1.9 Å. Moreover, most basic residues (blue sticks) and all basic residues (red sticks) are solvent exposed. All dipole moment values (Table S2) for the candidate structure of feature A2/B1 were calculated in YASARA using the AMBER03 force field. A total of seven acidic residues (D$^{24}$, E$^{52}$, E$^{53}$, E$^{63}$, D$^{68}$, E$^{74}$, and D$^{85}$) were deprotonated, while 24 of 25 basic residues (R$^{3}$, K$^{5}$, K$^{5}$, K$^{12}$, K$^{16}$, R$^{17}$, R$^{39}$, K$^{20}$, K$^{31}$, R$^{35}$, R$^{39}$, R$^{40}$, K$^{44}$, R$^{45}$, R$^{55}$, R$^{59}$, R$^{77}$, R$^{79}$, K$^{91}$, R$^{92}$, and K$^{95}$) were protonated, and the last basic residue was neutralized, leading to a combination of 25 structures having different charge distributions. The charge distribution dependence on the dipole is shown in Figure S$b$. A closer inspection revealed that the largest dipole moment value of 1149 D was observed for a structure with neutralized residue R$^{92}$ near the C-terminus.
whereas the smallest dipole moment value of 664 D was observed for a structure with neutralized residue K44 at the N-terminus. While we cannot exclude other potential candidates for similar CCS values, this example demonstrates that the charge distribution, particularly in structures having exposed acidic and basic residues, can lead to a wide distribution of the dipole alignment. These results are in agreement with the previous report that suggested dipole alignment under FAIMS experiments for molecular ions having dipole values above 300 D.40

■ CONCLUSIONS

In this work, DNA-free H4.1 structures have been addressed using FAIMS-TIMS-MS in parallel with MD simulations. The TIMS-TOF MS experiments under soft ion injection conditions showed that H4.1 adopts different structures ranging from C-like to P-like and subsequently to E-like conformations. The FAIMS-TIMS-MS results showed that native-like structures were retained throughout the FAIMS stage for the lowest charge states (6+ and 7+). However, P-like structures (z = 9+ to 11+) were sensitive to ion heating in the FAIMS stage. Moreover, no change in the CCS profiles were observed for the highest charge states by the addition of the FAIMS stage, allowing the assessment of the influence of the conformational diversity and charge distribution on the dipole alignment at high electric fields. The results from CIU experiments revealed heterogeneous structures for H4.1. Interestingly, “refolding after unfolding” upon activation was observed for the lowest charge state (z = 7+). The MD simulations revealed the native conformation of H4.1 with both N- and C-terminal domains folded toward the α-helical core. The gas-phase unfolding pathways showed that H4.1 exhibited major structural rearrangements associated with unfolding of the N- and C-terminal vicinities and disruption of inter-α-helical interactions, leading to the exposure of charge residues and increasing the distance of N- and C-terminal domains in higher charge states.

Complementary CCS and dipole calculations illustrated that charge distributions in an E-like structure, where basic and acidic residues are exposed (e.g., z = 17+), are sufficient to induce differences in the dipole alignment at high electric fields, particularly when their CCS profiles exhibit a narrow band distribution.

In conclusion, the data offer important insights into DNA-free H4.1 structures and differences when compared with H2A.1. These data can be further used as guidelines for structure-based drug design specifically for histones.

■ EXPERIMENTAL SECTION

Materials and Reagents. Recombinant human histone H4.1 (accession number: P62805) was purchased from EpiCypher (Durham, NC). All solvents used in this study were of analytical grade or better and purchased from Fisher Scientific (Pittsburgh, PA). Low-concentration Tuning Mix standard (G1969-85000) was purchased from Agilent Technologies (Santa Clara, CA) and used as received. Stock H4.1 was prepared in ultrapure water and extensively dialyzed against 10 mM ammonium acetate (NH₄Ac) buffer for ~36 h by changing buffer every ~12 h. Native H4.1 samples were prepared by diluting histone stock to 10 μM and 65 mM NH₄Ac solutions, whereas the denatured H4.1 sample was prepared at 10 μM in a methanol/water/formic acid (v/v 50:49:1) solution.

nESI-FAIMS-TIMS-TOF MS Instrument. Details regarding the operation of the nanoelectrospray ionization (nESI)-FAIMS-TIMS-TOF MS instrument can be found elsewhere.21 Briefly, the system comprised a “dome” FAIMS unit with a 2 mm annular gap (g)41–43 orthogonally mounted to a custom-built TIMS-TOF MS (Bruker Daltonics Inc., MA).44 The FAIMS system was controlled by a custom driver outputting a bisinusoidal waveform with a frequency of 1 MHz and an amplitude (dispersion voltage, DV) of 4 kV in the negative polarity mode applied to the inner electrode (GAA Custom Electronics, LLC). The compensation voltage (CV) was linearly scanned with a rate (Sr_FAIMS) of 1.0 V/min and expressed in the compensation field (E_c) by the equation E_c = CV/g.

A nESI emitter was made from a quartz capillary (o.d. = 1.0 mm) using a P2000 laser puller (Sutter Instruments Co., Novato, CA). The voltages for the nESI emitter and the curtain plate were set at 2.4 and 1.0 kV, respectively. N₂ gas was purified by a GC filter (Agilent, RSMN-4) and used to carry ions into the dome FAIMS. The carrier gas flow was controlled at 2 L/min by a digital flowmeter (MKS Instruments).

In TIMS, nitrogen (N₂) was used as a buffer gas at a constant velocity controlled by the pressure difference between the funnel entrance (P₁ = 2.6 mbar) and exit (P₂ = 1.1 mbar) at ca. 300 K.44 A radio frequency (rf) voltage of 250 Vpp at a frequency of 880 kHz was applied to TIMS electrodes. The reduced mobility (K₀) of ions depended on the gas flow velocity (ν_g), elution voltage (V_e), base voltage (V_out), and ramp time (t_ramp) and can be expressed by eq 1

\[
K_0 = \frac{V_e}{E} \approx \frac{A}{V_e - V_{\text{out}}} \tag{1}
\]

where A is a calibration constant determined using the known mobility of Tuning Mix ions m/z 622.029 (1/K₀ = 0.985 Vs/cm²), 922.009 (1/K₀ = 1.190 Vs/cm²), 1221.991 (1/K₀ = 1.382 Vs/cm²), and 1821.952 (1/K₀ = 1.729 Vs/cm²).25,45 The K₀ of an ion is correlated to the collision cross section (CCS, Ω) via the Mason–Schamp equation (eq 2)46

\[
\Omega = \frac{(18\pi)^{1/2}}{16} \left(\frac{z}{kT}\right)^{1/2} \left[\frac{1}{m_1} + \frac{1}{m_2}\right]^{1/2} \frac{1}{K_0} \frac{1}{N^8} \tag{2}
\]

where z is the charge of the ion, kB is the Boltzmann constant, T is the temperature, N⁸ is the bath gas number density, and m₁ and m₂ are the masses of the molecular ion and the bath gas, respectively.25,45

To avoid potential ion activation prior to mobility analysis, all experiments were carried out using soft ion transmission and trapping conditions.47,48 A voltage difference (∆V) between 10 to 20 V was maintained between the deflector (V_def) and the funnel entrance (V_fun), as well as between the funnel entrance and the TIMS analyzer (V_ramp). A V_def of −140 V, a V_fun of −160 V, a base voltage (V_out) of 60V, and a V_ramp of −180 to −20 V were used for all measurements. The TIMS scan rate (Sr_TIMS) was set at 1.6 V/ms.

The reference FAIMS-TOF MS spectra were collected on the TIMS transmission mode (without trapping), while the reference TIMS-TOF MS spectra were collected with the FAIMS unit detached. The 4D FAIMS-FAIMS-TIMS-TOF MS data were acquired by summing the TIMS-MS spectra (~100
spectra) every 0.9 V/cm over the $E_{\text{def}} = -15\text{\,V/cm}$ range. Data were processed using Data Analysis 5.0 (Bruker Daltonics Inc., Billerica, MA) and Microcal Origin 7.0. (OriginLab, Northampton, MA).

The CIU experiments were performed by varying $V_{\text{def}}$ with a 20 V interval. Voltages $V_{\text{fin}}$ and $V_{\text{tang}}$ were changed relative to $V_{\text{def}}$ to ensure a homogeneous axial electric field in all CIU experiments.29 The CIU data were analyzed using CIU Suite 2 software.50

**Classical Molecular Dynamics (CMD) Simulations.** CMD simulations were carried out as previously described.21 Simulations were performed in explicit solvent at 310 K using NAMD 2.14 (NVIDIA CUDA acceleration)53 with the modified CHARMM36 (CHARMM36m) force field for intrinsically disordered proteins.54 Long-range interactions were treated with the particle mesh Ewald method with a 12 Å nonbonded cutoff.55 To clarify the effects of water box thickness and nonbonded cutoff, CMD simulations of a system using a water box cutoff of 25 Å and a nonbonded cutoff of 8 Å were performed. The system was energy minimized for 500 000 steps using the conjugate gradient and the line search algorithm, followed by heating at 1 K/ps from 0 to 310 K. After reaching 310 K, a 100 ps equilibration run with an integration time step of 1 fs in the NVT ensemble was performed using Langevin dynamics to maintain the temperature at 310 K. The production run of 500 ns was performed in the NPT ensemble at 1 atm and 310 K with Langevin coupling and an integration time step of 2 fs. Trajectory files were saved every 0.1 ns for analysis.

**Simulated Annealing Molecular Dynamics (SAMD) Simulations.** SAMD simulations were used to generate gas-phase structures.21,28,56 The smallest radius of gyration ($R_{\text{gyr}} = 14.8$ Å) structure obtained from the CMD (conformation H, Figure 4b) was used as the starting structure. To ensure oversampling of total conformation spaces, SAMD simulations were performed for 220 cycles with the temperature ramping between T1 (300 K) and T2 (800 K): 50 ps relaxation at T1, 5 ps heating to T2, 50 ps relaxation at T2, and 2 ps cooling to T1, 5 ps minimization, and a 20 K increase between T1 and T2. An NVT ensemble with an integration time step of 2 fs was used. While this method generated most of the candidates (extended compact 7+ to unfolded 18+), the compact structures for 6+ to 9+ were obtained using a supervised NVT ensemble with boundary restraints to force the conversion into more compact structures followed by energy minimization.

**Steered Molecular Dynamics (SMD) Simulations.** Details regarding the gas-phase SMD simulations can be found elsewhere.21,28,57 The extended unfolding structures of H4.1 were generated by pulling the center of mass of the α-carbons of residue S at position 1 (S1) on the N-terminus at a rate of 0.25 Å/ns for 100 ns while anchoring residue G at position 102 (G102) on the C-terminus. The production run was performed in the NVT ensemble with an integration time step of 2 fs. The trajectory structures were saved at every 0.5 ps for further analysis. The reproducibility and reliability of SMD simulations were also clarified.

**Theoretical CCS Calculation and Cluster Analysis.** The theoretical CCS values of H4.1 in $N_2$ gas (CCS$_{\text{gas}}$) were calculated using the trajectory method (TM) integrated with IMoS v.1.10b software.58–61 The candidate structures of H4.1 are classified based on the comparison of root-mean-square deviation (RMSD) for all structures and classified into clusters.58,62 The statistical processing codes62 written in R (https://www.r-project.org) were used in this work. When candidates from the cluster families were determined, charge assignment was performed by scoring the relative solvent accessibility (RSA) of basic and acidic residues in H4.1.63 The dipole moments of individual candidates were computed with the AMBER03 force field using YASARA software (www.yasara.org). The electric dipole moment $\mathbf{p}$ in the unit of C·m is the length of vector $\mathbf{P}$, which is defined by eq 3

$$\mathbf{p} = \sum q_i (\mathbf{R}_i - \bar{C})$$

(3)

In this formula, $q_i$ is the charge of atom $i$, $\mathbf{R}$ is its Cartesian coordinate, and $C$ is the geometric center. Finally, all of the structural models were displayed using PyMOL (Schrödinger, Inc.).

**ASSOCIATED CONTENT**

Supporting Information
The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsomega.1c03744.

Complete FAIMS-TIMS analyses of native and denatured H4.1. CIU-TIMS fingerprints of all charge states; candidate structures of all charge states; table of all experimental and theoretical CCS values and a table of dipole values; and structural coordinates for in-solution and gas-phase candidate structures of H4.1 (PDF)

**AUTHOR INFORMATION**

**Corresponding Author**
Francisco Fernandez-Lima — Department of Chemistry and Biochemistry, Florida International University, Miami, Florida 33199, United States; Biomolecular Science Institute, Florida International University, Miami, Florida 33199, United States; orcid.org/0000-0002-1283-4390; Email: fernandf@fiu.edu

**Author**
Khoa N. Pham — Department of Chemistry and Biochemistry, Florida International University, Miami, Florida 33199, United States; orcid.org/0000-0001-7449-1226

Complete contact information is available at: https://pubs.acs.org/10.1021/acsomega.1c03744

**Notes**
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