CURVATURE PROPERTIES OF INTERIOR BLACK HOLE METRIC

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(Received 22 August 2019; accepted 5 December 2019)

A spacetime is a connected 4-dimensional semi-Riemannian manifold endowed with a metric tensor $g$ with signature $(-+++)$. The geometry of a spacetime is described by the tensor $g$ and the Ricci tensor $S$ of type $(0, 2)$ whereas the energy momentum tensor of type $(0, 2)$ describes the physical contents of the spacetime. Einstein’s field equations relate $g$, $S$ and the energy momentum tensor and describe the geometry and physical contents of the spacetime. By solving Einstein’s field equations for empty spacetime (i.e. $S = 0$) for a non-static spacetime metric, one can obtain the interior black hole solution, known as the interior black hole spacetime which infers that a remarkable change occurs in the nature of the spacetime, namely, the external spatial radial and temporal coordinates exchange their characters to temporal and spatial coordinates, respectively, and hence the interior black hole spacetime is a non-static one as the metric coefficients are time dependent. For the sake of mathematical generalizations, in the literature, there are many rigorous geometric structures constructed by imposing the restrictions to the curvature tensor of the space involving first order and second order covariant differentials of the curvature tensor. Hence a natural question arises that which geometric structures are admitted by the interior black hole metric. The main aim of this paper is to provide the answer of this question so that the geometric structures admitting by such a metric can be interpreted physically.
**Key words**: Einstein’s field equations; interior black hole metric; warped product metric; Tachibana tensor; quasi-Einstein manifold; 2-quasi-Einstein manifold; partially Einstein manifold; pseudosymmetric space; curvature condition of pseudosymmetry type.

**2010 Mathematics Subject Classification**: 53B20, 53B25, 53B50, 53C25, 53C40, 83C57.

1. **INTRODUCTION**

In the theory of general relativity one of the exciting predictions is that there may exist regions of the spacetime, where the gravity is so strong that nothing not even light, can ever escape. Such regions are known as black hole of the spacetime. It is well known that the most general spherically symmetric, static, vacuum, asymptotically flat exact solution to Einstein’s field equations is described by the Schwarzschild metric

\[ ds^2 = -\left(1 - \frac{2m}{z}\right)dt^2 + \left(1 - \frac{2m}{z}\right)^{-1}dz^2 + z^2(d\theta^2 + \sin^2\theta d\phi^2), \tag{1.1} \]

where \(z, \theta, \phi\) are spherical polar coordinates, \(t\) is the time measure by a clock at infinity and \(m = \frac{MG}{c^2}\), \(M\) being the mass of the central body, \(G\) being the gravitational constant and \(c\) being the velocity of light. Geometric properties of the Schwarzschild metric are presented among others in \([6, \text{Chapters 23-26}], [65, \text{Section 8}]\) and \([80, \text{Chapter 13}]\). We also refer to \([15, \text{Chapter 4.7}]\) for a short introduction to this subject. Curvature properties of pseudosymmetry type of this metric are given in \([53, 102]\).

On the sphere \(z = 2m\), the coefficient of \(dz^2\) in (1.1) tends to infinity and hence \(z = 2m\) is a singularity. The metric has also a singularity at \(z = 0\) as the coefficient at \(dt^2\) in (1.1) tends to infinity. We note that \(z = 0\) is the centre of the spherical mass distribution of the star. Since the metric coefficients are coordinate dependent, \(z = 0\) is a coordinate singularity, which implies that the inverse of the metric components \(g_{33}\) and \(g_{44}\) diverges even though at that point there is nothing such physical possibility. The geometry behaves very strangely when \(z = 2m\) and it is called the Schwarzschild radius. The spherical surface associated with the Schwarzschild radius \(z = 2m\) is null and corresponds to the black hole’s “event horizon”, where a freely falling particle can approach to the surface \(z = 2m\) but never cross it and hence the inward falling particle need infinite time to reach the surface of the sphere \(z = 2m\), which was first pointed out by Oppenheimer and Snyder \([81]\). We note that inside the Schwarzschild radius, \(z\) and \(t\) coordinates change their role in the sense that the \(t\) coordinate becomes spacelike and \(z\) coordinate becomes timelike. It is believed that the gravitational collapse of a compact body results in a singularity hidden beyond an event horizon. If the singularity were visible to the exterior region one would have a naked singularity which would open the realm...
for wild speculations [56]. This entails to Penrose’s cosmic censorship conjecture [82] which states that all physically reasonable spacetimes are globally hyperbolic, forbidding the existence of naked singularities, and only allowing singularities to be hidden behind event horizon. The realization that black hole could actually exists prompted a renewed interest in their mathematical properties and the last three decades have seen some remarkable developments in this respect. For details about the black hole in cosmology and astrophysics we refer the article of Carr [10] and also references therein.

The empty annular region of spacetime for a spherical star inside its Schwarzschild radius \(2m\) and outside its physical radius \(a\), \(a < 2m\), that is for a black hole the spacetime geometry is characterized by the spherical symmetric non-static line element

\[
ds^2 = -B(z, t)dt^2 + A(z, t)dz^2 + t^2(d\theta^2 + \sin^2 \theta d\phi^2),
\]

(1.2)

where the coefficient functions can be obtained by solving Einstein’s field equations for empty space-time \(S = 0\). The solution takes the form

\[
ds^2 = -\left(\frac{2\xi}{t} - 1\right)^{-1}dt^2 + \left(\frac{2\xi}{t} - 1\right)dz^2 + t^2(d\theta^2 + \sin^2 \theta d\phi^2).
\]

(1.3)

Thus (1.3) represents metric for interior black hole [56, 70], where \(\xi\) may be determined from a direct confrontation with the exterior Schwarzschild solution. The metric (1.3) is the interior black hole solution which represents the empty spacetime in the exterior region \(z > a\) of a black hole and it is valid for \(a\) such that \(2m < z < a\). For physical significance and cosmological interpretation of the interior black hole solution we refer the reader [56] and also references therein. However in the interior black hole solution, a remarkable change occurs in the nature of spacetime namely the external spatial radial and temporal coordinates exchange their characters to temporal and spatial coordinates, respectively and hence the interior black hole solution is represented by a non-static spacetime as its metric coefficients are time dependent.

The nature of a space is completely known by its curvature which can be explicitly determined by the metric of that space. In the literature of differential geometry there are several kinds of generalizations of various geometrical structures constructed by giving the curvature restrictions involving first and second order covariant derivatives. The notion of local symmetry is a generalization of the manifold of constant curvature and the study was initiated by Cartan in 1926 with full classification of such a space [11, 12]. Also a full classification of locally symmetric semi-Riemannian space is given by Cahen and Parker [8, 9]. During the last eight decades the process of generalization of locally symmetric spaces have been carried out by many authors around the globe in different directions, for instance, recurrent manifold by Walker [118], 2-recurrent manifold by Lichnerowicz [79],
quasi-generalized recurrent manifold by Shaikh and Roy [104], hyper-generalized recurrent manifold by Shaikh and Patra [103], weakly generalized recurrent manifold by Shaikh and Roy [85, 105], semisymmetric manifold by Cartan [13], pseudosymmetric manifold by Chaki [14], pseudosymmetric manifold by Deszcz [28, 38], Ricci-pseudosymmetric manifold by Deszcz and Hotloś [41], weakly symmetric manifold by TÁMassy and Binh [113], weakly symmetric manifold by Selberg [84]. We mention that pseudosymmetry by Chaki and Deszcz are different and also weak symmetry of Selberg and TÁMassy and Binh are different.

We consider the semi-Riemannian manifold \((M, g)\), \(\dim M = 4\), equipped with the interior black hole metric given in (1.3). Then this manifold is a interior black hole spacetime. The main subject of this paper is to investigate the geometric structures admitting by the interior black hole spacetime. The paper is organized as follows. In Section 2 we present definitions of some special tensors. In Section 3 we present basical facts on pseudosymmetric manifolds (in the sense of Deszcz [114]). In Section 4 we deduce the curvature properties of interior black hole metric and found that interior black hole spacetime is a pseudosymmetric manifold.

Finally, in the last section (Appendix) we present the local components of the considered tensors of the metric (1.3). We also mention that we have made all the calculations by a programme in Wolfram Mathematica.

2. SOME SPECIAL TENSORS

Let \((M, g)\), \(n = \dim M \geq 3\), be a connected smooth semi-Riemannian manifold with Levi-Civita connection \(\nabla\) and semi-Riemannian metric \(g\). For \((0, 2)\)-tensors \(A\) and \(B\) on \(M\) we define their Kulkarni-Nomizu product \(A \wedge B\) by (see, e.g., [36, 60, 94, 95, 96, 100, 106])

\[
(A \wedge B)(X_1, X_2, X, Y) = A(X_1, Y)B(X_2, X) + A(X_2, X)B(X_1, Y) - A(X_1, X)B(X_2, Y) - A(X_2, Y)B(X_1, X).
\]

We define the endomorphisms \(X \wedge_A Y\), \(\mathcal{R}(X, Y)\), \(\mathcal{C}(X, Y)\), \(\mathcal{P}(X, Y)\), \(\mathcal{W}(X, Y)\) and \(\mathcal{K}(X, Y)\) by [18, 30, 34, 42, 43, 60, 62]

\[
(X \wedge_A Y)Z = A(Y, Z)X - A(X, Z)Y,
\]

\[
\mathcal{R}(X, Y)Z = [\nabla_X, \nabla_Y]Z - \nabla_{[X,Y]}Z,
\]

\[
\mathcal{C}(X, Y) = \mathcal{R}(X, Y) - \frac{1}{n-2}(X \wedge g \mathcal{L}Y + \mathcal{L}X \wedge g Y - \frac{r}{n-1}X \wedge g Y),
\]

\[
\mathcal{P}(X, Y) = \mathcal{R}(X, Y) - \frac{1}{n-1}X \wedge_S Y,
\]

\[
\mathcal{W}(X, Y) = \mathcal{R}(X, Y) - \frac{1}{n-1}X \wedge_S Y,
\]

\[
\mathcal{K}(X, Y) = \mathcal{R}(X, Y) - \frac{1}{n-1}X \wedge_S Y.
\]
\[ W(X, Y) = R(X, Y) - \frac{r}{n(n-1)} X \wedge g Y, \]
\[ K(X, Y) = R(X, Y) - \frac{1}{n-2} (X \wedge g LCY + LX \wedge g Y), \]

respectively, where \( A \) is a \((0,2)\)-tensor on \( M \), \( X, Y, Z \in \chi(M) \), \( \chi(M) \) being the Lie algebra of smooth vector fields on \( M \). The Ricci tensor \( S \), the Ricci operator \( \mathcal{L} \) and the scalar curvature \( r \) are defined by \( S(X, Y) = \text{tr} \{ Z \mapsto \mathcal{R}(Z, X)Y \} \), \( g(\mathcal{L}X, Y) = S(X, Y) \) and \( r = \text{tr} \mathcal{L} \), respectively. We define the tensor \( G \), the Riemannian-Christoffel curvature tensor \( R \), the Weyl conformal curvature tensor \( C \), the projective curvature tensor \( P \), the concircular curvature tensor \( W \) and the conharmonic curvature tensor \( K \) of \((M, g)\), by [18, 30, 34, 42, 43, 60, 62]

\[
G(X_1, \ldots, X_4) = g((X_1 \wedge g X_2)X_3, X_4) = \frac{1}{2}(g \wedge g)(X_1, \ldots, X_4),
\]
\[
R(X_1, \ldots, X_4) = g(\mathcal{R}(X_1, X_2)X_3, X_4),
\]
\[
C(X_1, \ldots, X_4) = g(C(X_1, X_2)X_3, X_4)
\]
\[
\quad = (R - \frac{1}{n-2} g \wedge S + \frac{r}{(n-1)(n-2)} G)(X_1, \ldots, X_4),
\]
\[
P(X_1, \ldots, X_4) = g(P(X_1, X_2)X_3, X_4)
\]
\[
\quad = R(X_1, \ldots, X_4) - \frac{1}{n-1}(g(X_1, X_4)S(X_2, X_3) - g(X_2, X_4)S(X_1, X_3)),
\]
\[
W(X_1, \ldots, X_4) = g(W(X_1, X_2)X_3, X_4) = (R - \frac{r}{n(n-1)} G)(X_1, \ldots, X_4),
\]
\[
K(X_1, \ldots, X_4) = g(K(X_1, X_2)X_3, X_4) = (R - \frac{1}{n-2} g \wedge S)(X_1, \ldots, X_4)
\]
\[
\quad = (C - \frac{r}{(n-1)(n-2)} G)(X_1, \ldots, X_4),
\]

respectively. For an \((0, k)\)-tensor \( T \), \( k \geq 1 \), and a symmetric \((0,2)\)-tensor \( A \) we define the \((0, k+2)\)-tensor \( Q(A, T) \) [94, 95, 98, 99] by

\[
Q(A, T)(X_1, \ldots, X_k; X, Y) = ((X \wedge_A Y) \cdot T)(X_1, \ldots, X_k)
\]
\[
\quad = -T((X \wedge_A Y)X_1, X_2, \ldots, X_k) - \cdots - T(X_1, \ldots, X_{k-1}, (X \wedge_A Y)X_k).
\]

The tensor \( Q(A, T) \) is called the Tachibana tensor of the tensors \( A \) and \( T \), or shortly the Tachibana tensor [32, 33]. It is obvious that the tensor \( Q(g, G) \) vanishes identically on any semi-Riemannian manifold. Therefore we have \( Q(g, R) = Q(g, W) \) and \( Q(g, C) = Q(g, K) \). For an endomorphism \( \mathcal{D}(X, Y) \) we define the \((0, 4)\)-tensor \( D \) by

\[
D(X_1, \ldots, X_4) = g(\mathcal{D}(X_1, X_2)X_3, X_4).
\]
Now for an \((0, k)\)-tensor \(T, k \geq 1\), and an endomorphism \(D(X, Y)\) we define the \((0, k+2)\)-tensor \(D \cdot T\) [94, 95, 98, 99] by

\[
(D \cdot T)(X_1, \ldots, X_k; X, Y) = (D(X, Y) \cdot T)(X_1, \ldots, X_k)
\]

\[
= -T(D(X, Y)X_1, X_2, \ldots, X_k) - \cdots - T(X_1, \ldots, X_{k-1}, D(X, Y)X_k).
\]

Setting in the above formulas \(D(X, Y) = \mathcal{R}(X, Y), C(X, Y), \mathcal{P}(X, Y), W(X, Y), K(X, Y), T = R, S, C, P, W, K\) and \(A = g\) or \(S\), we obtain the tensors: \(R \cdot R, R \cdot S, R \cdot C, R \cdot P, R \cdot W, R \cdot K, C \cdot R, C \cdot S, C \cdot C, C \cdot P, C \cdot W, C \cdot K, P \cdot R, P \cdot S, P \cdot C, P \cdot P, P \cdot W, P \cdot K, W \cdot R, W \cdot S, W \cdot C, W \cdot P, W \cdot W, W \cdot K, K \cdot R, K \cdot S, K \cdot C, K \cdot P, K \cdot W, K \cdot K, Q(g, R), Q(g, S), Q(g, C), Q(g, P), Q(g, W), Q(g, K), Q(S, R), Q(S, C), Q(S, P), Q(S, W), Q(S, K).

Using the above presented definitions we can prove that the following identities hold on any semi-Riemannian manifold \((M, g), n \geq 4\), (cf. [32, Proposition 1.1]): \(R \cdot K = R \cdot C\) and

\[
K \cdot S = C \cdot S - \frac{r}{(n-1)(n-2)}Q(g, S),
\]

\[
K \cdot R = C \cdot R - \frac{r}{(n-1)(n-2)}Q(g, R),
\]

\[
K \cdot K = C \cdot C - \frac{r}{(n-1)(n-2)}Q(g, C).
\]

Moreover, we also have \(R \cdot W = R \cdot R, R \cdot C = R \cdot K, C \cdot R = C \cdot W, C \cdot C = C \cdot K, W \cdot R = W \cdot W, W \cdot C = W \cdot K, K \cdot R = K \cdot W, K \cdot C = K \cdot K, P \cdot R = P \cdot W\) and \(P \cdot C = P \cdot K\).

### 3. Pseudosymmetry Type Manifolds

A semi-Riemannian manifold \((M, g), n \geq 3\), is said to be an Einstein manifold [4] if at every point of \(M\) its Ricci tensor \(S\) is proportional to the metric tensor \(g\), i.e. \(S = \frac{r}{n}g\) on \(M\). In particular, if \(S\) vanishes on \(M\) then it is called Ricci flat. We denote by \(U_S\) the set of all points of \(M\) at which \(S\) is not proportional to \(g\), i.e. \(U_S = \{x \in M : S - \frac{r}{n}g \neq 0\ \text{at}\ x\}\).

As a generalization of Einstein manifold, the notion of quasi-Einstein manifold arose during the study of exact solution of Einstein field equations as well as during the investigation of quasi-umbilical hypersurfaces of conformally flat spaces. For instance, FLRW spacetimes are quasi-Einstein spacetimes. The semi-Riemannian manifold \((M, g), n \geq 3\), is said to be a quasi-Einstein manifold if rank \((S - \alpha g) = 1\) on \(U_S \subset M\), where \(\alpha\) is some function on this set (see, e.g., [16, 18, 29, 34, 47, ...]}
48, 52, 57, 62, 90, 91, 92, 107, 108]). In particular, if rank $S = 1$ on $U_S$ then $(M, g)$ is called Ricci-simple [21]. For instance, the Gödel spacetime is a Ricci-simple manifold (see, e.g., [45]). The semi-Riemannian manifold $(M, g)$, $n \geq 3$, is said to be a 2-quasi-Einstein manifold if rank $(S - \alpha g) \leq 2$ on $U_S \subset M$ and rank $(S - \alpha g) = 2$ on some open non-empty subset of $U_S$, where $\alpha$ is some function on $U_S$ (see, e.g., [35, 36, 37, 88, 95, 101]). Such manifolds also are called generalized quasi-Einstein manifolds, cf. [57, 58] and references therein. It is easy to check that every non-Einstein warped product manifold with an 1-dimensional base and a semi-Riemannian Einsteinian $(n-1)$-dimensional fibre is a quasi-Einstein manifold. Similarly, it is easy to check that every non-quasi-Einstein warped product manifold with an 2-dimensional base and a semi-Riemannian Einsteinian $(n-2)$-dimensional fibre is a 2-quasi-Einstein manifold.

An extension of the class of Einstein semi-Riemannian manifolds $(M, g)$, $n \geq 3$, also form Ricci-symmetric manifolds [11-13], i.e. manifolds with parallel Ricci tensor ($\nabla S = 0$). We note that the scalar curvature of every Ricci-symmetric manifold is constant. In [64] Gray introduced two classes of manifolds lying between the class of Ricci-symmetric manifolds and the class of manifolds of constant scalar curvature, viz., the class $A$ is the class of manifolds which are cyclic Ricci parallel $((\nabla_X S)(Y, Z) + (\nabla_Y S)(Z, X) + (\nabla_Z S)(X, Y) = 0)$ and the class $B$ is the class with Codazzi type Ricci tensor $((\nabla_X S)(Y, Z) = (\nabla_Y S)(X, Z))$. Existence of both classes are given in [86] (see also [45]). Codazzi type Ricci tensor was extensively studied by various authors [3, 4, 7, 22, 23, 24, 59, 109]. Another important subclass of the class of Ricci-symmetric manifolds form locally symmetric manifolds [11-13], i.e. manifolds for which we have $\nabla R = 0$. This implies the following integrability condition $\mathcal{R}(X, Y) \cdot R = 0$, in short

$$R \cdot R = 0.$$ (3.1)

A semi-Riemannian manifold $(M, g)$, $n \geq 3$, is called semisymmetric [13] if (3.1) holds on $M$ and a full classification of such manifolds, in the Riemannian case, is given by Szabó [110-112]. Further, a semi-Riemannian manifold $(M, g)$, $n \geq 3$, is said to be pseudosymmetric (or, in the sense of Deszcz) [17, 28, 38, 40, 66, 67, 68, 89, 98, 99, 114, 116, 117] if the tensors $R \cdot R$ and $Q(g, R)$ are linearly dependent at every point of $M$. This is equivalent to

$$R \cdot R = L_R Q(g, R)$$ (3.2)

on $U_R = \{ x \in M : R - \frac{r}{n(n-1)} G \neq 0 \text{ at } x \}$, where $L_R$ is a function on $U_R$. Pseudosymmetric manifolds (in the sense of Deszcz [114]) are also called Deszcz symmetric spaces (see, e.g., [17, 68, 115, 117]). A pseudosymmetric manifold is called a pseudosymmetric space of constant type if the function $L_R$ is constant [75, 76]. We mention that a geometrical interpretation of (3.2), in the Riemannian case, is given in [67].
We note that pseudosymmetric tensors arose during the study of semisymmetric totally umbilical submanifolds in manifolds admitting semisymmetric generalized curvature tensors [1, 25, 27]. For example, every totally umbilical submanifold of a semisymmetric manifold, with parallel mean curvature vector, is pseudosymmetric [1, 2]. The systematic study on pseudosymmetric manifolds was initiated in [1]. We refer to [40, 68] for a wider presentation related to the last statement. We mention that [38] is the first publication, in which a semi-Riemannian manifold satisfying (3.2) was called the pseudosymmetric manifold.

The Schwarzschild spacetime, the Kottler spacetime, the Reissner-Nordström spacetime and the Reissner-Nordström-de Sitter spacetime satisfy (3.2) with non-zero function $L_R$ [53] (see also [39, 66]). We also refer to [18, 20, 49, 50, 51, 74] for further results on pseudosymmetric spacetimes. For instance, a family of curvature conditions satisfied by the Reissner-Nordström-de Sitter spacetime was determined in [74]. The Schwarzschild spacetime was discovered in 1916 by Schwarzschild and independently by Droste, during their study on solutions of Einstein’s equations, see, e.g., [6, Section 23.3] and [83] and references therein. It seems that the Schwarzschild spacetime, the Reissner-Nordström spacetime, as well as some Friedmann-Lemaître-Robertson-Walker (FLRW) spacetimes are the “oldest” examples of a non-semisymmetric pseudosymmetric warped product manifolds (cf. [36, 40]).

Pseudosymmetric manifolds form a subclass of the class of Ricci pseudosymmetric manifolds. A semi-Riemannian manifold $(M, g)$, $n \geq 3$, is said to be Ricci-pseudosymmetric ([28, 41]) if the tensors $R \cdot S$ and $Q(g, S)$ are linearly dependent at every point of $M$. This is equivalent to

$$R \cdot S = L_S Q(g, S)$$  \hspace{1cm} (3.3)

on $U_S$, where $L_S$ is a function on this set. Ricci-pseudosymmetric manifolds are also called Ricci Deszcz symmetric spaces (see, e.g., [117]). A Ricci-pseudosymmetric manifold is called a Ricci-pseudosymmetric manifold of constant type if the function $L_S$ is constant [61]. It is obvious that (3.2) implies (3.3). The converse statement is not true, provided that $n \geq 4$, (see, e.g., [33]). However, (3.2) and (3.3) are equivalent on every 3-dimensional semi-Riemannian manifold. The conditions (3.2) and (3.3) are equivalent on every 4-dimensional warped products [26]. The conditions (3.2) and (3.3) are also equivalent on hypersurfaces isometrically immersed in 5-dimensional semi-Riemannian space of constant curvature [19]. It is known that every warped product manifold with an 1-dimensional base and a semi-Riemannian Einsteinian $(n-1)$-dimensional fibre is a Ricci-pseudosymmetric manifold [33, 41, 52]. For further results on Ricci-pseudosymmetric manifolds we refer to [33]. We mention that a geometrical interpretation of (3.3), in the Riemannian case, is given in [72].
We denote by $\mathcal{U}_C$ the set of all points of a semi-Riemannian manifold $(M, g)$, $n \geq 4$, at which $C \neq 0$. We note that $\mathcal{U}_S \cup \mathcal{U}_C = \mathcal{U}_R$, see, e.g., [31].

A semi-Riemannian manifold $(M, g)$, $n \geq 4$, is said to be a manifold with pseudosymmetric Weyl tensor [26, 28, 31, 40, 54, 55] if the tensors $C \cdot C$ and $Q(g, C)$ are linearly dependent at every point of $M$. This is equivalent to

$$C \cdot C = L_C Q(g, C)$$

on $\mathcal{U}_C$, where $L_C$ is a function on this set. Every warped product manifold with an 2-dimensional base and a 2-dimensional fibre is a manifold with pseudosymmetric Weyl tensor [26, Theorem 2]. Recently in [36] it was proved that this statement is also true when the fibre is an $(n-2)$-dimensional space of constant curvature, $n \geq 4$. Thus in particular, the 4-dimensional spacetime with the metric (1.2), as well as the 5-dimensional spacetime with the metric (3.7) are 2-quasi-Einstein manifolds with pseudosymmetric Weyl tensor. It may be mentioned that the Gödel spacetime satisfies (3.4) (see, [45]). We refer to [45, 46, 87, 88, 97, 101, 102, 107] for examples of various pseudosymmetric type structures.

As it was stated in [55, Theorem 3.1], if $(M, g)$, $n \geq 4$, is a pseudosymmetric manifold with pseudosymmetric Weyl tensor then

$$Q(S - \alpha g, C - \beta G) = 0$$

on $\mathcal{U}_C$, where $\alpha$ and $\beta$ are some functions on this set. Moreover, from (3.5) it follows that at all points of $\mathcal{U}_S \cap \mathcal{U}_C$, at which rank $(S - \alpha g) > 1$, we have (cf. [55, Theorem 3.2])

$$R = L_1 S \wedge S + L_2 g \wedge S + L_3 g \wedge g,$$

where $L_1$, $L_2$ and $L_3$ are some functions on this set. We refer to [33, 49, 50, 51, 63, 73, 74] for results on manifolds satisfying (3.6). The manifold satisfying (3.6) is said to be Roter type manifold. Roter type manifolds are also called Roter spaces.

Some comments on pseudosymmetric manifolds (also called Deszcz symmetric spaces), as well as Roter spaces, are given in [17, Section 1]: "From a geometric point of view, the Deszcz symmetric spaces may well be considered to be the simplest Riemannian manifolds next to the real space forms." and "From an algebraic point of view, Roter spaces may well be considered to be the simplest Riemannian manifolds next to the real space forms." For further comments we refer to [117]. Recently Roter spaces admitting geodesic mappings were studied in [44].
4. Geometric Structures Admitting by Interior Black Hole Spacetime

Let \((B, \tilde{g})\) and \((F, \tilde{g})\) be semi-Riemannian manifolds of dimension \(p \geq 1\) and \(n - p \geq 1\), respectively, covered by the coordinate charts \(\{U; x^a\}\) and \(\{V; y^a\}\), respectively. Let \(f\) be a smooth positive function on \(B\). The warped product \(M = B \times_f F\) is the product manifold \(B \times F\) furnished with the metric \(g = \pi^*(g_B) + (f \circ \pi)^*(g_F)\), where \(\pi\) and \(\sigma\) are the projections of \(B \times F\) onto \(B\) and \(F\), respectively [77, 78]. The manifold \(B\) is called the base of \(M = B \times F\), and \(F\) the fiber. We mention that for the warped product manifold \(B \times_f F\), the metric can also be considered \([5, 15]\) as \(g = \pi^*(g_B) + (f \circ \pi)^2g^*(g_F)\). However, throughout the paper we will consider the former warped product metric but not later.

Let \(\{\mathcal{U} \times \tilde{V}; x^1, ..., x^p, x^{p+1} = y^1, ..., y^{n-p}\}\) be a product chart for \(B \times F\). The local components of the metric \(g = \tilde{g} \times_f \tilde{g}\) with respect to this chart are given by the following:

\[
g_{hk} = \tilde{g}_{ab}\text{ if } h = a \text{ and } k = b, \quad g_{hk} = f\tilde{g}_{\alpha\beta}\text{ if } h = \alpha \text{ and } k = \beta \text{ and } g_{hk} = 0, \quad \text{otherwise,}
\]

where \(a, b, c, \ldots \in \{1, \ldots, p\}\), \(\alpha, \beta, \gamma, \ldots \in \{p + 1, \ldots, n\}\) and \(h, i, j, k, l, m \in \{1, 2, \ldots, n\}\). We will mark by bars (resp., by tildes) objects formed from \(\tilde{g}\) (resp. \(\tilde{g}\)).

The local components \(\Gamma^h_{jk}\) of the Levi-Civita connection \(\nabla\) of \(B \times_f F\) are given by the following:

\[
\Gamma^a_{bc} = \tilde{\Gamma}^a_{bc}, \quad \Gamma^a_{\beta\gamma} = \tilde{\Gamma}^a_{\beta\gamma}, \quad \Gamma^a_{ab} = \Gamma^a_{ab} = 0, \quad (4.1)
\]

\[
\Gamma^a_{\beta\gamma} = -\frac{1}{2}\tilde{g}^{ab}f_a\tilde{g}_{\gamma\beta}, \quad \Gamma^a_{a\beta} = \frac{1}{2f}f_a\delta^a_{\beta}, \quad f_a = \partial_a f = \frac{\partial f}{\partial x^a}, \quad (4.2)
\]

The local components

\[
R_{hijk} = g_{hi}R^l_{ijk} = g_{hi}(\partial_k \Gamma^l_{ij} - \partial_j \Gamma^l_{ik} + \Gamma^m_{ij} \Gamma^l_{mk} - \Gamma^m_{ik} \Gamma^l_{mj}), \quad \partial_k = \frac{\partial}{\partial x^k},
\]

of the Riemann-Christoffel curvature tensor \(R\) and the local components \(S_{ij}\) of the Ricci tensor \(S\) of the warped product \(\mathcal{B} \times_f \tilde{F}\) which may not vanish identically are the following:

\[
R_{abcd} = \tilde{R}_{abcd}, \quad R_{a\beta\gamma\delta} = -\frac{1}{2}T_{ab}\tilde{g}_{\alpha\beta}, \quad R_{a\beta\gamma\delta} = f\tilde{R}_{a\beta\gamma\delta} = \frac{\Delta_1 f}{4}\tilde{G}_{a\beta\gamma\delta}, \quad (4.3)
\]

\[
S_{ab} = \tilde{S}_{ab} - \frac{n-p}{2f}T_{ab}, \quad S_{a\beta} = \tilde{S}_{a\beta} - \frac{1}{2}\left(\text{tr}(T) + \frac{n-p-1}{2f}\Delta_1 f\right)\tilde{g}_{a\beta}, \quad (4.4)
\]

\[
T_{ab} = \overline{\nabla}_b f_a - \frac{1}{2f}f_a f_b, \quad \text{tr}(T) = \tilde{g}^{ab}T_{ab}, \quad \Delta_1 f = \tilde{g}^{ab}f_a f_b, \quad (4.5)
\]

and \(T\) is the \((0, 2)\)-tensor with the local components \(T_{ab}\). The scalar curvature \(r\) of \(\mathcal{B} \times_f \tilde{F}\) satisfies the following:

\[
r = \overline{r} + \frac{\overline{r}}{f} - \frac{n-p}{f}\left(\text{tr}(T) + \frac{n-p-1}{4f}\Delta_1 f\right), \quad (4.6)
\]
Curvature properties of interior black hole metric

For further details about warped products, we refer to [80]. Warped product pseudosymmetric and Ricci pseudosymmetric manifolds are studied among others in [18, 26, 38, 50, 51, 54]. Also we refer to [93] for the weakly symmetric and weakly Ricci symmetric warped product manifolds.

Let \((M, g)\) be the manifold with the metric \(g\) defined by (1.3). Using the above presented formulas we can compute the local components of tensors formed by the metric tensor defined by (1.3) (see Section 5). We set \(\xi = \frac{d\xi}{dt}\) and \(\dot{\xi} = \frac{d\dot{\xi}}{dt}\). We can check that \(S = (r/4)g\) at all points of \(M\) at which \(\xi = 2\dot{\xi}/t\) and \(\text{rank}(S + (2\dot{\xi}/t^2)g) = 2\) at remains points of \(M\), i.e. on \(U_S \subset M\). Thus \((M, g)\) is a 2-quasi-Einstein manifold. We can also check that \(S^2 = (r/2)(S - (2\dot{\xi}/t^3)g)\), and in a consequence, \(S^2 - (\text{tr}(L^2)/4)g = (r/2)(S - (r/4)g)\) on \(U_S \subset M\), where \(S^2\) is a \((0,2)\)-tensor with the local components \(S^2_{ij} = S_{ik}g^{kl}S_{lj}\) and \(\text{tr}(L^2) = g^{kl}g^2_{kl}\). Thus the considered manifold is a partially Einstein manifold.

We recall that a semi-Riemannian manifold \((M, g), n \geq 3\), is said to be partially Einstein manifold [15, p. 20] if \(S^2 = \alpha S + \beta g\) on \(U_S \subset M\), where \(\alpha\) and \(\beta\) are some functions on this set. It is easy to verify that every quasi-Einstein manifold is partially Einstein (see, e.g., [62, Introduction]). It is obvious that the converse statement is not true. We also have [115, 117]: A conformally flat semi-Riemannian manifold of dimension \(n \geq 4\) is a Deszcz symmetric space if and only if it is partially Einstein.

We present now results related to the interior black hole metric (1.3).

Theorem 4.1 — Interior black hole metric (1.3) satisfies the following:

(i) \(R \cdot Z = L_1Q(g, Z), \quad L_1 = \frac{\xi - t\dot{\xi}}{t^3}\),
(ii) \(C \cdot Z = L_2Q(g, Z), \quad L_2 = \frac{6\dot{\xi} - 4t\dot{\xi}^2 + 2t^2\dot{\xi}}{6t^3}\),
(iii) \(W \cdot Z = L_2Q(g, Z),
(iv) \(K \cdot Z = L_3Q(g, Z), \quad L_3 = \frac{2\dot{\xi} + t^2\dot{\xi}}{2t^3}\),
(v) \(P \cdot S = L_1Q(g, S),
(vi) \(P \cdot Z = L_1Q(g, Z) - \frac{1}{3}Q(S, Z),
where \(Z\) is any one of \(R, S, C, W, K\) and \(P\).

From above theorem it follows that (a) if \(\xi = C_1 t\) then \(R \cdot Z = 0\) and \(P \cdot S = 0\), (b) if \(\xi = C_1 t^2 + C_2 t^3\) then \(C \cdot Z = 0\) and \(W \cdot Z = 0\), and (c) if \(\xi = \sqrt{t} \left( C_2 \cos \left( \frac{\sqrt{7}}{2} \log t \right) + C_1 \sin \left( \frac{\sqrt{7}}{2} \log t \right) \right)\) then \(K \cdot Z = 0\), where \(C_1\) and \(C_2\) are some constants. Further, we have

Theorem 4.2 — Interior black hole metric (1.3) satisfies the following:

(i) \(C \cdot K = W \cdot K, \quad C \cdot K = W \cdot C, \quad C \cdot C = W \cdot C,
(ii) \(W \cdot K = W \cdot C, \quad C \cdot K = C \cdot C \quad ((ii) follows from (i)),\)
(iii) \( W \cdot K = C \cdot C \) ((iii) follows from (i) and (ii)),

(iv) \( C \cdot W = W \cdot R \),

(v) \( R \cdot S = P \cdot S \), \( C \cdot S = W \cdot S \),

(vi) \( L_3 R \cdot K = L_1 K \cdot C \),

(vii) \( R \cdot W - W \cdot R = L_5 Q(g, R) \), \( L_5 = -\frac{2\xi + t\xi}{6t^2} = \frac{c}{12}, \)

(viii) \( C \cdot K - K \cdot C = L_6 Q(g, C) \), \( L_6 = -\frac{2\xi + t\xi}{3t^2} = \frac{c}{6} \),

(ix) \( C \cdot R - Q(S, C) = L_7 Q(g, C) \), \( L_7 = \frac{3t\xi + 4t^2 \xi}{2t^3} \),

(x) \( R \cdot R - Q(S, R) = L_8 Q(g, C) \), \( L_8 = \frac{6\xi^2 - 4t\xi \xi - 2t^2 \xi^2 + 4t^2 \xi \xi}{t^3(6\xi - 4t \xi + t^2 \xi)} \),

(xi) \( L_6 L_9 R \cdot W + L_1 L_{10} Q(S, W) = L_1 L_{11} Q(S, R) \), \( L_9 = -3t(6\xi(\xi - t\xi) + t(t\xi + 3\xi)\xi) \), \( L_{10} = 6(-t^2 \xi^2 + 2t\xi(t\xi - \xi) + 3\xi^2) \), \( L_{11} = t^2(-t^2 \xi^2 + 2t\xi^2 + 2t\xi \xi) + 6t(\xi \xi - 4\xi) + 18\xi^2 \),

(xii) \( L_3 R \cdot K + L_{12} K \cdot R + L_3 Q(S, K) = 0 \), \( L_{12} = -\frac{2\xi + 4t\xi + t^2 \xi}{2t^3} \),

(xiii) \( L_3 L_{13} R \cdot K - L_2 L_2 K \cdot R + L_1 L_2 L_3 Q(S, R) = 0 \), \( L_{13} = \frac{3\xi^2 - 2t\xi \xi - 2t^2 \xi \xi + 2t^2 \xi \xi}{3t^6} \),

(xiv) \( L_2 C \cdot W - L_7 W \cdot C = L_2 Q(S, C) \),

(xv) \( L_2 L_{14} C \cdot W - \frac{1}{18t^6} L_{11} W \cdot C = L_2 Q(S, W) \), \( L_{14} = \frac{3t\xi - 5t\xi - t^2 \xi}{3t^6} \),

(xvi) \( L_{15} C \cdot K + L_2^2 Q(S, K) = L_2 L_{16} Q(S, C) \), \( L_{15} = \frac{4t^2 \xi + 4t^2 \xi + t^2 \xi}{3t^6} \), \( L_{16} = \frac{2t^2 \xi - 2t^2 \xi}{2t^3} \),

(xvii) \( 18t^6 L_{13} L_{11} W \cdot K - L_2^2 L_1 K \cdot W + L_2^2 L_3 Q(S, W) = 0 \),

(xviii) \( L_1 L_3 W \cdot K + L_2 L_{12} K \cdot W + L_2 L_3 Q(S, K) = 0 \),

(xix) Roter type condition with \( R = \frac{\phi}{2} S \wedge S + \mu g \wedge S + \eta g \), where

\[
\phi = -\frac{6t\xi - 4t^2 \xi + t^3 \xi}{(t\xi - \xi)^2}, \quad \mu = -\frac{6t\xi - 6t\xi^2 + 3t\xi \xi + t^2 \xi \xi}{t(t\xi - \xi)^2}, \quad \eta = \frac{2(4t\xi^2 - 4t\xi^3 + 2t\xi \xi \xi + t^2 \xi \xi)}{t^3(t\xi - \xi)^2}.
\]

From above theorem it follows that (a) if \( \xi = C_1 t \) then \( R \cdot K = K \cdot C \), (b) if \( \xi = -\frac{C_1}{t} + C_2 \) then \( R \cdot W = W \cdot R \) and \( C \cdot K = K \cdot C \), (c) if \( \xi = t \left( C_2 \cos \left( \frac{\sqrt{23}}{4} \log t \right) + C_1 \sin \left( \frac{\sqrt{23}}{4} \log t \right) \right) \) then \( C \cdot R = Q(S, C) \), where \( C_1 \) and \( C_2 \) are some constants.

In [50] it was proved that any Roter type manifold satisfies the relation (x) with \( L_8 = L_1 + \phi^{-1} \mu = (n - 2) \phi^{-1} (\mu^2 - \phi \eta) \), \( n = 4 \), and the converse is also true as follows from (x) and (xix).

Let \( (M, g) \) be the manifold with the interior black hole metric in 5-dimension is given by [69-71]

\[
 ds^2 = - \left( \frac{2\xi}{t^2} - 1 \right)^{-1} dt^2 + \left( \frac{2\xi}{t^2} - 1 \right) dz^2 + t^2 (d\theta^2 + \sin^2 \theta d\phi^2 + \sin^2 \theta \sin^2 \phi d\psi^2),
\]

where \( \xi \) is the function of time. Similarly as above, we can obtain curvature properties of the interior black hole metric (3.7). In particular, we again set \( \dot{\xi} = \frac{dt}{dt} \) and \( \ddot{\xi} = \frac{d^2 t}{dt^2} \). We can check that \( S = (r/5)g \) at all points of \( M \) at which \( \ddot{\xi} = 3\dot{\xi}/t \) and rank\( (S + (2\dot{\xi}/t^3)g) = 2 \) at remains points of \( M \), i.e. on \( U_S \subset M \). Thus \( (M, g) \) is a 2-quasi-Einstein manifold. We can also check that \( S^2 - (tr(L^2)/5)g = \)
Further, we have the following curvature properties of the considered metric (3.7).

**Theorem 4.3** — 5-dimensional interior black hole metric (3.7) satisfies the following:

(i) \( R \cdot Z = N_1 Q(g, Z), \quad N_1 = \frac{2t - t^2}{t^3} \),

(ii) \( C \cdot Z = N_2 Q(g, Z), \quad N_2 = \frac{12t - 6t^2 + t^2 \xi}{6t^3} \),

(iii) \( W \cdot Z = N_3 Q(g, Z), \quad N_3 = \frac{20t - 8t^2 + t^2 \xi}{10t^3} \),

(iv) \( K \cdot Z = N_4 Q(g, Z), \quad N_4 = \frac{6t - 2t^2 \xi + t^2 \xi}{3t^4} \),

(v) \( P \cdot S = N_1 Q(g, S), \)

(vi) \( P \cdot Z = N_1 Q(g, Z) - \frac{1}{4} Q(S, Z) \),

where \( Z \) is any one of \( R, S, C, W, K \) and \( P \).

From above theorem it follows that (a) if \( \xi = C_1 t^2 \) then \( R \cdot Z = 0 \) and \( P \cdot S = 0 \), (b) if \( \xi = C_1 t^3 + C_2 t^4 \) then \( C \cdot Z = 0 \), (c) if \( \xi = C_1 t^5 + C_2 t^4 \) then \( W \cdot Z = 0 \), and (d) if \( \xi = t^{\frac{3}{2}} \left( C_2 \cos \left( \frac{\sqrt{3} \pi}{2} \log t \right) + C_1 \sin \left( \frac{\sqrt{3} \pi}{2} \log t \right) \right) \) then \( K \cdot Z = 0 \), where \( C_1 \) and \( C_2 \) are some constants.

**Theorem 4.4** — 5-dimensional interior black hole metric (3.7) satisfies the following:

(i) \( R \cdot S = P \cdot S \),

(ii) \( R \cdot W - W \cdot R = N_5 Q(g, R), \quad N_5 = \frac{-2t + t^2}{10t^3} \),

(iii) \( C \cdot K - K \cdot C = N_6 Q(g, C), \quad N_6 = \frac{-2t + t^2}{6t^3} \),

(iv) \( C \cdot R - Q(S, C) = N_7 Q(g, C), \quad N_7 = \frac{4t^2 + t^3}{2t^3} \),

(v) \( R \cdot R - Q(S, R) = N_8 Q(g, C), \quad N_8 = \frac{3(8t^2 - 4t^2 \xi - t^2 \xi^2 + 2t^2 \xi^3)}{t^3 (12t - 6t^2 + t^2 \xi)} \),

(vi) \( N_5 N_9 R \cdot W + N_1 N_9 Q(S, R) + N_1 N_8 Q(S, W) = 0, \quad N_9 = \frac{-3t^2 + 4t^3 \xi + 6t^2 + t^2 \xi}{2t^4} \),

\( N_{10} = \frac{-80t^2 + 64t^2 \xi - 2t^2 \xi^2 - 4t^2 \xi (2t + t^3) + t^4 \xi^2}{20t^8} \),

(vii) \( N_4 R \cdot K + N_{11} K \cdot R + N_4 Q(S, K) = 0, \quad N_{11} = \frac{-6t^2 + 6t^3 \xi + t^2 \xi}{3t^4} \),

(viii) \( N_4 N_8 R \cdot K - N_1^2 N_2 K \cdot R + N_1 N_2 N_4 Q(S, R) = 0, \)

(ix) \( N_3 C \cdot W - N_7 W \cdot C = N_3 Q(S, C), \)

(x) \( N_3 N_{12} C \cdot W + N_{10} W \cdot C = N_3 N_2 Q(S, W), \quad N_{12} = \frac{20t^2 - 16t^2 \xi - 3t^2 \xi}{10t^3} \),

(xi) \( N_{13} C \cdot K + N_2 N_{11} Q(S, C) + N_2^2 Q(S, K) = 0, \quad N_{13} = \frac{6t^2 + 5t^2 \xi + t^2 \xi^2}{6t^6} \),

(xii) \( N_4 N_{10} W \cdot K + N_{12} N_{14} K \cdot W = N_4 N_{14} Q(S, W), \quad N_{14} = \frac{240t^2 - 216t^2 \xi + 48t^2 \xi^2 + 32t^2 \xi (14t^2 \xi + t^2 \xi^2)}{6t^6} \),

(xiii) \( N_1 N_4 W \cdot K + N_3 N_{11} K \cdot W + N_3 N_4 Q(S, K) = 0 \),

\( ((r/2) + (\xi/2))(S - (r/5)g) \) on \( U_S \subset M \). Thus the considered manifold is a partially Einstein manifold.
(xiv) **Roter type condition** with \( R = \frac{\phi}{2} S \wedge S + \mu g \wedge S + \eta G \), where

\[
\phi = -\frac{t^2(12\xi + t(\xi \dot{\xi} - 6\dot{\xi}))}{(3\xi - t\dot{\xi})^2}, \quad \mu = -\frac{3\dot{\xi}(4\xi - 3t\dot{\xi}) + t(4\xi + t\dot{\xi})\ddot{\xi}}{t(3\xi - t\dot{\xi})^2},
\]

\[
\eta = \frac{-2(-6t\xi^3 + \xi(9\dot{\xi}^2 + 2t\ddot{\xi} + t^2\dddot{\xi}))}{t^4(3\xi - t\dot{\xi})^2}.
\]

From above theorem it follows that (a) if \( \xi = -\frac{C_1}{t} + C_2 \) then \( R \cdot W = W \cdot R \) and \( C \cdot K = K \cdot C \), (b) if \( \xi = \sqrt{t} \left( C_2 \cos \left( \frac{\sqrt{15}}{2} \log t \right) + C_1 \sin \left( \frac{\sqrt{15}}{2} \log t \right) \right) \) then \( C \cdot R = Q(S, C) \), and (c) if \( \xi = \frac{C_2(3+10t^2C_1)}{t^2} \) then \( R \cdot R = Q(S, R) \), where \( C_1 \) and \( C_2 \) are some constants.

However, we note that for the 5-dimensional interior black hole metric the following tensors are non-zero tensors

(i) \( C \cdot K - W \cdot K \), (ii) \( C \cdot K - W \cdot C \), (iii) \( C \cdot C - W \cdot C \), (iv) \( C \cdot W - W \cdot R \), (v) \( C \cdot S - W \cdot S \) and (vi) \( R \cdot K - K \cdot C \).

In [50] it was proved that any Roter type manifold satisfies the relation (v) with \( N_8 = N_1 + \phi^{-1} \mu = (n - 2)\phi^{-1}(\mu^2 - \phi\eta) \), \( n = 5 \), and hence from (xiv), it follows that the converse of the result is also true.

We can check that the 4-dimensional spacetime with the metric (1.3) and the 5-dimensional spacetime with the metric (3.7) are non-quasi Einstein manifolds. Furthermore from the considerations presented in Section 3 and Theorems 4.1 (i) and 4.4 (ii) it follow that those spacetimes satisfy (3.6) and hence are Roter type spacetimes.

We also note that interior black hole metrics (1.3) and (3.7) do not admit any one of the following structures: Ricci semisymmetric, quasi-Einstein, Codazzi type Ricci tensor, cyclic Ricci symmetric, Chaki pseudo symmetric, Chaki pseudo Ricci symmetric, weakly symmetric, weakly Ricci symmetric, hyper generalized recurrent, weakly generalized recurrent, quasi generalized recurrent, as well as any pseudosymmetric type structure defined by \( R \cdot Z = LQ(S, Z) \), \( C \cdot Z = LQ(S, Z) \), \( W \cdot Z = LQ(S, Z) \), \( K \cdot Z = LQ(S, Z) \), \( P \cdot Z = LQ(S, Z) \), and \( P \cdot Z = LQ(g, Z) \), respectively, where \( L \) is any smooth function and \( Z \) is any one of the tensors \( R, C, W, K \) and \( P \). It can also be mentioned that both interior black hole metrics does not realize any one of the generalized Einstein metric condition (i) \( R \cdot C - C \cdot R = L_1Q(g, R) \), (ii) \( R \cdot C - C \cdot R = L_2Q(g, C) \), (iii) \( R \cdot C - C \cdot R = L_3Q(S, R) \) and (iv) \( R \cdot C - C \cdot R = L_4Q(S, C) \). We mention that a survey on semi-Riemannian manifolds satisfying the last four conditions is given in [33].
5. Appendix

Part I: From (1.3) and (4.1)-(4.6), the local components of the Christoffel symbols of second kind, the curvature tensor and the Ricci tensor (upto symmetry) which may not vanish identically are the following:

\[ \Gamma^1_{11} = -\frac{\xi - t\dot{\xi}}{t^2 - 2t\xi}, \quad \Gamma^3_{13} = \Gamma^4_{14} = \frac{1}{t}, \quad \Gamma^1_{22} = -\frac{(t - 2\xi)(t\dot{\xi} - \xi)}{t^3}, \]

\[ \Gamma^1_{33} = 2\xi - t, \quad \Gamma^4_{34} = \cot \theta, \quad \Gamma^1_{44} = -(t - 2\xi) \sin^2 \theta, \quad \Gamma^3_{44} = -\sin \theta \cos \theta, \]

\[ R_{1212} = -\frac{t^2\xi - 2t\xi + 2\xi}{t^3}, \quad R_{1313} = \frac{t\xi - \xi}{t - 2\xi}, \quad R_{1414} = \frac{\sin^2 \theta(t\dot{\xi} - \xi)}{t - 2\xi}, \]

\[ R_{2323} = -\frac{(t - 2\xi)(t\dot{\xi} - \xi)}{t^2}, \quad R_{2424} = -\frac{(t - 2\xi) \sin^2 \theta(t\dot{\xi} - \xi)}{t^2}, \quad R_{3434} = 2t\xi \sin^2 \theta, \]

\[ S_{11} = -\frac{\ddot{\xi}}{t - 2\xi} = -\frac{\ddot{\xi}}{t} g_{11}, \quad S_{22} = \frac{(t - 2\xi)\dddot{\xi}}{t^2} = -\frac{\dddot{\xi}}{t} g_{22}, \]

\[ S_{33} = -2\ddot{\xi} = -\frac{2\ddot{\xi}}{t^2} g_{33}, \quad S_{44} = -2\ddot{\xi} \sin^2 \theta = -\frac{2\ddot{\xi}}{t^2} g_{44}, \quad r = -\frac{2(\ddot{\xi} + t\dddot{\xi})}{t^2}, \]

where \( \dot{\xi} \) denotes the differentiation of \( \xi \) with respect to \( t \) and \( \dddot{\xi} \) the second order differentiation of \( \xi \) with respect to \( t \). The local components of the Weyl conformal curvature tensor (upto symmetry) which may not vanish identically are given by

\[ C_{1212} = -\frac{6\xi - t(4\dot{\xi} - t\ddot{\xi})}{3t^3}, \quad C_{1313} = -\frac{6\xi - t(4\dot{\xi} - t\ddot{\xi})}{6(t - 2\xi)}, \quad C_{1414} = -\frac{\sin^2 \theta[6\xi - t(4\dot{\xi} - t\ddot{\xi})]}{6(t - 2\xi)}, \]

\[ C_{2323} = \frac{(t - 2\xi)[6\xi - t(4\dot{\xi} - t\ddot{\xi})]}{6(t - 2\xi)}, \quad C_{2424} = \frac{(t - 2\xi) \sin^2 \theta[6\xi - t(4\dot{\xi} - t\ddot{\xi})]}{6(t - 2\xi)}, \]

\[ C_{3434} = -\frac{1}{3} t \sin^2 \theta[6\xi - t(4\dot{\xi} - t\ddot{\xi})]. \]

The local components of the covariant derivatives of curvature tensor and Ricci tensor (upto symmetry) which may not vanish identically are given by:

\[ R_{1212,1} = -\frac{t^3\xi^3 + 3t^2\dddot{\xi} - 6t\dot{\xi} + 6\xi}{t^4}, \quad R_{1223,3} = \frac{(t - 2\xi)(t(t\dddot{\xi} - 3\dot{\xi}) + 3\xi)}{t^3} = -R_{2323,1}, \]

\[ R_{1224,4} = \frac{(t - 2\xi) \sin^2 \theta(t(t\dddot{\xi} - 3\dot{\xi}) + 3\xi)}{t^3} = -R_{2424,1}, \quad R_{1313,1} = \frac{t^2\dddot{\xi} - 3t\dot{\xi} + 3\xi}{t^2 - 2t\xi}, \]

\[ R_{1334,4} = \sin^2 \theta(3\xi - t\ddot{\xi}) = -R_{1434,3} = -\frac{1}{2} R_{3434,1}, \quad R_{1414,1} = \frac{\sin^2 \theta(t(t\dddot{\xi} - 3\dot{\xi}) + 3\xi)}{t(t - 2\xi)}, \]
\[
S_{11,1} = \frac{\dot{\xi} - t \xi^3}{t^3 - 2t \xi}, \quad S_{13,3} = \frac{2\dot{\xi} - \ddot{\xi}}{t} = \frac{1}{2} S_{33,1},
\]
\[
S_{22,1} = \frac{(t - 2\xi)(t \xi^3 - \ddot{\xi})}{t^3 - 2t \xi}, \quad S_{14,4} = \frac{\sin^2(2\dot{\xi} - t \ddot{\xi})}{t} = \frac{1}{2} S_{44,1}.
\]

In terms of local coordinate system, the local components \( Q(A, T)_{i_1 i_2 \ldots i_k} \) of the Tachibana tensor \( Q(A, T) \) of an \((0, 2)\)-tensor \( A \) and an \((0, k)\)-tensor \( T \) are given by

\[
Q(A, T)_{i_1 i_2 \ldots i_k} = A_{i_1} u T_{i_2 \ldots i_k} + A_{i_2} u T_{i_1 \ldots i_k} + \cdots + A_{i_k} u T_{i_1 \ldots i_{k-1}} - A_{i_k} v T_{i_1 \ldots i_{k-1}} - \cdots - A_{i_2} v T_{i_1 i_k} - A_{i_1} v T_{i_2 \ldots i_k}.
\]

In particular, for a symmetric \((0, 2)\)-tensor \( A \) and a generalized curvature tensor \( T \), we have

\[
Q(A, T)_{hijklm} = Q(A, T)_{jkhilm} = -Q(A, T)_{ihjklm} = -Q(A, T)_{hjkml},
\]
\[
Q(A, T)_{hijklm} = Q(A, T)_{ijhklm} = Q(A, T)_{hklijm},
\]
\[
Q(A, T)_{hijklm} + Q(A, T)_{jklmh} + Q(A, T)_{lmhijk} = 0.
\]

If \( A \) and \( B \) are symmetric \((0, 2)\)-tensors then

\[
Q(A, B)_{hijk} = Q(A, B)_{ihjk} = -Q(A, B)_{hikj}.
\]

If \( D(X, Y) = R(X, Y) \) then (2.1) yields

\[
(R \cdot T)_{i_1 i_2 \ldots i_k} = -g^{pq}(T_{p i_1 \ldots i_k} R_{uvq} + T_{i_p \ldots i_k} R_{uvqi} + \cdots + T_{i_1 i_2 \ldots p} R_{uvqi_k}),
\]
where \( g^{pq}, R_{hijk} \) and \( T_{i_1 i_2 \ldots i_k} \) are the local components of the tensors \( g^{-1}, R \) and \( T \), respectively. Similarly in terms of local coordinate system we can write the components of \( C \cdot T, P \cdot T, W \cdot T \) and \( K \cdot T \). Moreover, if \( B_{hijk} \) and \( T_{hijk} \) are the local components of generalized curvature tensors \( B \) and \( T \) then the local components of the \((0, 6)\)-tensor \( B \cdot T \) are following

\[
(B \cdot T)_{hijklm} = -g^{pq}(T_{p i_1 i_2 i_3} B_{lmq} + T_{hp jk} B_{lmq} + T_{hijk} B_{lmq} + T_{hijp} B_{lmq}) = g^{pq}(T_{p j} B_{lmiq} - T_{ph jk} B_{lmq} + T_{pkhi} B_{lmq} - T_{pj hi} B_{lmq}).
\]

We have

\[
(B \cdot T)_{hijklm} = (B \cdot T)_{jkhilm} = -(B \cdot T)_{ihjklm} = -(B \cdot T)_{hjkml},
\]
\[
(B \cdot T)_{hijklm} + (B \cdot T)_{ijhklm} + (B \cdot T)_{jhiklm} = 0.
\]
For the tensors $R \cdot R$, $R \cdot S$, $R \cdot C$ and $R \cdot P$ we have the following relations:

$$
(R \cdot R)_{12313} = \frac{(t^2 - \xi)(t(t^2 - 3\xi) + 3\xi)}{t^4} = -(R \cdot R)_{121323},
$$

$$
(R \cdot R)_{143413} = \frac{\sin^2 \theta(t^2 - 3\xi)(t^2 - \xi)}{t(t - 2\xi)} = -(R \cdot R)_{133414},
$$

$$
(R \cdot R)_{122414} = \frac{\sin^2 \theta(t^2 - \xi)(t(t^2 - 3\xi) + 3\xi)}{t^4} = (R \cdot R)_{121424},
$$

$$
-(R \cdot R)_{243423} = \frac{(t^2 - 2\xi)\sin^2 \theta(t^2 - 3\xi)(t^2 - \xi)}{t^3} = (R \cdot R)_{233424};
$$

$$
(R \cdot S)_{1313} = \frac{(\xi - t^2)(t^2 - 2\xi)}{t^2(t - 2\xi)},
(R \cdot S)_{1414} = \frac{\sin^2 \theta(t^2 - \xi)(t\xi - 2\xi)}{t^2(t - 2\xi)},
$$

$$
(R \cdot S)_{2323} = \frac{(t^2 - 2\xi)(t\xi - 2\xi)}{t^4},
(R \cdot S)_{2424} = \frac{(t^2 - 2\xi)\sin^2 \theta(t\xi - \xi)(t\xi - 2\xi)}{t^3};
$$

$$
(R \cdot C)_{12313} = \frac{(t^2 - \xi)(t(t^2 - 3\xi) + 6\xi)}{2t^4} = -(R \cdot C)_{121323},
$$

$$
(R \cdot C)_{143413} = \frac{\sin^2 \theta(\xi - t^2)(t(t^2 - 3\xi) + 6\xi)}{2t(t - 2\xi)} = -(R \cdot C)_{133414},
$$

$$
(R \cdot C)_{122414} = \frac{\sin^2 \theta(t^2 - \xi)(t(t^2 - 3\xi) + 6\xi)}{2t^4} = -(R \cdot C)_{121424},
$$

$$
(R \cdot C)_{243423} = \frac{(t^2 - 2\xi)\sin^2 \theta(t^2 - \xi)(t(t^2 - 3\xi) + 6\xi)}{2t^3} = -(R \cdot C)_{233424};
$$

$$
(R \cdot P)_{12313} = -(R \cdot P)_{131223} = -(R \cdot P)_{232113} = -(R \cdot P)_{121323} = \frac{(t^2 - \xi)(t(t^2 - 3\xi) + 3\xi)}{t^4},
$$

$$
(R \cdot P)_{123213} = -(R \cdot P)_{132123} = -(R \cdot P)_{231213} = -(R \cdot P)_{123123} = \frac{(\xi - t^2)(t(2t^2 - 7\xi) + 9\xi)}{3t^4},
$$

$$
(R \cdot P)_{131113} = \frac{(t^2 - \xi)(t\xi - 2\xi)}{3t(t - 2\xi)^2},
(R \cdot P)_{133313} = \frac{(\xi - t^2)(t\xi - 2\xi)}{3(t - 2\xi)};
$$

$$
(R \cdot P)_{143413} = -(R \cdot P)_{341314} = (R \cdot P)_{143413} = -(R \cdot P)_{341413} = \frac{\sin^2 \theta(\xi - t^2)(t(t\xi - 5\xi) + 9\xi)}{3t(t - 2\xi)},
$$

$$
(R \cdot P)_{144313} = (R \cdot P)_{344113} = (R \cdot P)_{133414} = -(R \cdot P)_{343114} = \frac{-\sin^2 \theta(t^2 - 3\xi)(t^2 - \xi)}{t(t - 2\xi)},
$$
\[(R \cdot P)_{122414} = -(R \cdot P)_{242114} = -(R \cdot P)_{121424} = -(R \cdot P)_{141224} = \frac{\sin^2 \theta(t\ddot{\xi} - \xi) \left(t(t\ddot{\xi} - 3\xi) + 3\xi \right)}{t^4},\]

\[(R \cdot P)_{124214} = -(R \cdot P)_{241214} = -(R \cdot P)_{124214} = -(R \cdot P)_{141214} = \frac{\sin^2 \theta(\xi - t\ddot{\xi}) \left(t(2t\ddot{\xi} - 7\xi) + 9\xi \right)}{3t^4},\]

\[\begin{align*}
(R \cdot P)_{141114} &= \frac{\sin^2 \theta(t\ddot{\xi} - \xi)(t\ddot{\xi} - 2\dot{\xi})}{3t(t - 2\xi)^2}, & (R \cdot P)_{144414} &= \frac{\sin^4 \theta(\xi - t\ddot{\xi})(t\ddot{\xi} - 2\dot{\xi})}{3(t - 2\xi)}, \\
(R \cdot P)_{232223} &= \frac{(t - 2\xi)^2(t\ddot{\xi} - \xi)(t\ddot{\xi} - 2\dot{\xi})}{3t^5}, & (R \cdot P)_{233323} &= \frac{(t - 2\xi)(t\ddot{\xi} - \xi)(t\ddot{\xi} - 2\dot{\xi})}{3t^2}, \\
(R \cdot P)_{244423} &= (R \cdot P)_{344223} = (R \cdot P)_{234324} = -t(2\xi) \sin^2 \theta(t\ddot{\xi} - \xi) \left(t(t\ddot{\xi} - 5\dot{\xi}) + 9\xi \right), \quad (R \cdot P)_{242224} = \frac{(t - 2\xi)^2 \sin^2 \theta(t\ddot{\xi} - \xi)(t\ddot{\xi} - 2\dot{\xi})}{3t^5},
\end{align*}\]

For the tensors $Q(g, R)$, $Q(g, S)$, $Q(g, C)$ and $Q(g, P)$ we have the following relations:

\[Q(g, R)_{122313} = -t\ddot{\xi} + 3\dot{\xi} - \frac{3\xi}{t} = -Q(g, R)_{121323},\]

\[Q(g, R)_{143413} = \frac{t^2 \sin^2 \theta(3\xi - t\dot{\xi})}{t - 2\xi} = -Q(g, R)_{133414},\]

\[Q(g, R)_{121424} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi \right)}{t} = -Q(g, R)_{122414},\]

\[Q(g, R)_{243423} = (t - 2\xi) \sin^2 \theta(t\ddot{\xi} - 3\dot{\xi}) = -Q(g, R)_{233424};\]

\[Q(g, S)_{131313} = \frac{t(t\ddot{\xi} - 2\dot{\xi})}{t - 2\xi}, \quad Q(g, S)_{141414} = \frac{\sin^2 \theta(t\ddot{\xi} - 2\dot{\xi})}{t - 2\xi},\]

\[Q(g, S)_{2323} = \frac{(t - 2\xi)(t\ddot{\xi} - 2\dot{\xi})}{t}, \quad Q(g, S)_{2424} = \frac{(t - 2\xi) \sin^2 \theta(t\ddot{\xi} - 2\dot{\xi})}{t};\]

\[Q(g, C)_{122313} = -\frac{1}{2} t\ddot{\xi} + 2\dot{\xi} - \frac{3\xi}{t} = -Q(g, C)_{121323},\]

\[Q(g, C)_{143413} = \frac{t^2 \sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi \right)}{2(t - 2\xi)} = -Q(g, C)_{133414},\]

\[Q(g, C)_{121424} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi \right)}{2t} = -Q(g, C)_{122414},\]
\[ Q(g, C)_{23324} = \frac{1}{2} (t - 2\xi) \sin^2 \theta \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi \right) = -Q(g, C)_{243423}; \]

\[ Q(g, P)_{131223} = Q(g, P)_{232113} = Q(g, P)_{121323} = -Q(g, P)_{122313} = t\ddot{\xi} - 3\dot{\xi} + \frac{3\xi}{t}, \]

\[ Q(g, P)_{123213} = -Q(g, P)_{132123} = -Q(g, P)_{231213} = -Q(g, P)_{123123} = \frac{2}{3} t\ddot{\xi} - \frac{7\dot{\xi}}{3} + \frac{3\xi}{t}, \]

\[ Q(g, P)_{131113} = -\frac{t^2(t\ddot{\xi} - 2\dot{\xi})}{3(t - 2\xi)^2}, \quad Q(g, P)_{133313} = \frac{t^3(t\ddot{\xi} - 2\dot{\xi})}{3(t - 2\xi)}, \]

\[ Q(g, P)_{143413} = Q(g, P)_{134314} = Q(g, P)_{341413} = -Q(g, P)_{341314} = \frac{t^2 \sin^2 \theta \left( t(t\ddot{\xi} - 5\dot{\xi}) + 9\xi \right)}{3(t - 2\xi)}, \]

\[ Q(g, P)_{144313} = Q(g, P)_{134413} = Q(g, P)_{344113} = -Q(g, P)_{344314} = \frac{t^2 \sin^2 \theta (t\ddot{\xi} - 3\dot{\xi})}{t - 2\xi}, \]

\[ Q(g, P)_{141224} = Q(g, P)_{242114} = Q(g, P)_{124124} = -Q(g, P)_{122414} = \frac{\sin^2 \theta \left( t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi \right)}{t}, \]

\[ Q(g, P)_{141114} = -\frac{t^2 \sin^2 \theta (t\ddot{\xi} - 2\dot{\xi})}{3(t - 2\xi)^2}, \quad Q(g, P)_{144414} = \frac{t^3 \sin^4 \theta (t\ddot{\xi} - 2\dot{\xi})}{3(t - 2\xi)}, \]

\[ Q(g, P)_{232223} = -\frac{(t - 2\xi)^2(t\ddot{\xi} - 2\dot{\xi})}{3t^2}, \quad Q(g, P)_{233323} = -\frac{1}{3} t(t - 2\xi)(t\ddot{\xi} - 2\dot{\xi}), \]

\[ Q(g, P)_{243423} = Q(g, P)_{342423} = Q(g, P)_{234324} = -Q(g, P)_{343224} = \frac{1}{3} (t - 2\xi) \sin^2 \theta \left( t(t\ddot{\xi} - 5\dot{\xi}) + 9\xi \right), \]

\[ Q(g, P)_{244323} = Q(g, P)_{344223} = Q(g, P)_{234224} = -Q(g, P)_{343224} = -(t - 2\xi) \sin^2 \theta (t\ddot{\xi} - 3\dot{\xi}), \]

\[ Q(g, P)_{242224} = -\frac{(t - 2\xi)^2 \sin^2 \theta (t\ddot{\xi} - 2\dot{\xi})}{3t^2}, \quad Q(g, P)_{244424} = -\frac{1}{3} t(t - 2\xi) \sin^4 \theta (t\ddot{\xi} - 2\dot{\xi}). \]

**Part II.** For the tensors \( C \cdot R, C \cdot S, C \cdot C \) and \( C \cdot P \) we have the following relations:

\[ (C \cdot R)_{121323} = \frac{t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi}{6t^4} \left( t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi \right) = -(C \cdot R)_{122313}, \]

\[ (C \cdot R)_{133414} = \frac{\sin^2 \theta (t\ddot{\xi} - 3\dot{\xi})}{6t(t - 2\xi)} \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi \right) = -(C \cdot R)_{143413}, \]

\[ (C \cdot R)_{121424} = \frac{\sin^2 \theta \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi \right)}{6t^4} \left( t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi \right) = -(C \cdot R)_{122414}, \]
(C \cdot R)_{243423} = \frac{(t - 2\xi) \sin^2 \theta (t\ddot{\xi} - 3\xi) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{6t^3} = -(C \cdot R)_{233424};

(C \cdot S)_{1313} = \frac{(t\ddot{\xi} - 2\dot{\xi}) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{6t^2(t - 2\xi)},

(C \cdot S)_{1414} = \frac{\sin^2 \theta (t\ddot{\xi} - 2\dot{\xi}) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{6t^2(t - 2\xi)},

(C \cdot S)_{2323} = \frac{(t - 2\xi)(t\ddot{\xi} - 2\dot{\xi}) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{6t^4},

(C \cdot S)_{2424} = \frac{(t - 2\xi) \sin^2 \theta (t\ddot{\xi} - 2\dot{\xi}) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{6t^4};

(C \cdot C)_{121323} = \frac{\left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)^2}{12t^4} = -(C \cdot C)_{122313},

(C \cdot C)_{143413} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)^2}{12t(t - 2\xi)} = -(C \cdot C)_{133414},

(C \cdot C)_{121424} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)^2}{12t^4} = -(C \cdot C)_{122414},

(C \cdot C)_{233424} = \frac{(t - 2\xi) \sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)^2}{12t^3} = -(C \cdot C)_{243423};

(C \cdot P)_{131223} = (C \cdot P)_{231213} = (C \cdot P)_{121323} = \frac{(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi) \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right)}{6t^4},

(C \cdot P)_{132123} = (C \cdot P)_{231213} = (C \cdot P)_{123123} = \frac{(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi) \left(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\xi\right)}{18t^4},

(C \cdot P)_{131113} = \frac{\left(t\ddot{\xi} - 2\dot{\xi}\right) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{18t(t - 2\xi)^2}, \quad (C \cdot P)_{133313} = \frac{(t\ddot{\xi} - 2\dot{\xi}) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{18(t - 2\xi)},

(C \cdot P)_{143413} = (C \cdot P)_{341413} = (C \cdot P)_{134314} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 5\dot{\xi}) + 9\xi\right) \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{18t(t - 2\xi)},
\[
(C \cdot P)_{144313} = (C \cdot P)_{344113} = (C \cdot P)_{133414} = -(C \cdot P)_{343114} = \frac{\sin^2 \theta (t \ddot{t} - 3 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{6t(t-2\dddot{t})},
\]
\[
(C \cdot P)_{141224} = (C \cdot P)_{242114} = (C \cdot P)_{121424} = -(C \cdot P)_{124214} = \frac{\sin^2 \theta (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t}) (t(t \dddot{t} - 3 \dddot{t}) + 3 \dddot{t})}{6t^4},
\]
\[
(C \cdot P)_{142124} = (C \cdot P)_{241214} = (C \cdot P)_{124124} = -(C \cdot P)_{124223} = \frac{\sin^2 \theta (t(t \dddot{t} - 2 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18t^6},
\]
\[
(C \cdot P)_{144114} = \frac{-\sin^2 \theta(t \dddot{t} - 2 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18(t-2 \dddot{t})},
\]
\[
(C \cdot P)_{232223} = -(t-2 \dddot{t})^2(t \dddot{t} - 2 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18t^5},
\]
\[
(C \cdot P)_{233323} = -(t-2 \dddot{t})(t \dddot{t} - 2 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18t^2},
\]
\[
(C \cdot P)_{243423} = (C \cdot P)_{342423} = (C \cdot P)_{234324}
\]
\[
= -(C \cdot P)_{342324} = -(t-2 \dddot{t}) \sin^2 \theta (t(t \dddot{t} - 5 \dddot{t}) + 9 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18t^3},
\]
\[
(C \cdot P)_{244323} = (C \cdot P)_{344223} = (C \cdot P)_{234424}
\]
\[
= -(C \cdot P)_{343224} = -(t-2 \dddot{t}) \sin^2 \theta(t \dddot{t} - 3 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{6t^3},
\]
\[
(C \cdot P)_{242224} = -(t-2 \dddot{t})^2 \sin^2 \theta(t \dddot{t} - 2 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18t^5},
\]
\[
(C \cdot P)_{244424} = -(t-2 \dddot{t}) \sin^4 \theta(t \dddot{t} - 2 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{18t^2}.
\]

For the tensors \( W \cdot R, W \cdot S, W \cdot C \) and \( W \cdot P \) we have the following relations:

\[
(W \cdot R)_{121323} = \frac{(t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t}) (t(t \dddot{t} - 3 \dddot{t}) + 3 \dddot{t})}{6t^4} = -(W \cdot R)_{122313},
\]
\[
(W \cdot R)_{133414} = \frac{\sin^2 \theta(t \dddot{t} - 3 \dddot{t}) (t(t \dddot{t} - 4 \dddot{t}) + 6 \dddot{t})}{6t(t-2 \dddot{t})} = -(W \cdot R)_{143413}.
\]
\[ (W \cdot R)_{121424} = \frac{\sin^2 \theta \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi} \right) \left( t(t\ddot{\xi} - 3\dot{\xi}) + 3\dot{\xi} \right)}{6t^4} = -(W \cdot R)_{224114}, \]

\[ (W \cdot R)_{243423} = \frac{(t - 2\xi) \sin^2 \theta (t\ddot{\xi} - 3\dot{\xi}) \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi} \right)}{6t^3} = -(W \cdot R)_{233424}; \]

\[ (W \cdot S)_{1313} = \frac{(t\ddot{\xi} - 2\dot{\xi}) \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi} \right)}{6t^2(t - 2\xi)}, \]

\[ (W \cdot S)_{1414} = \frac{\sin^2 \theta (t\ddot{\xi} - 2\dot{\xi})(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi}}{6t^2(t - 2\xi)}, \]

\[ (W \cdot S)_{2323} = -\frac{(t - 2\xi)(t\ddot{\xi} - 2\dot{\xi})(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})}{6t^4}, \]

\[ (W \cdot S)_{2424} = -\frac{(t - 2\xi) \sin^2 \theta (t\ddot{\xi} - 2\dot{\xi})(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})}{6t^4}; \]

\[ (W \cdot C)_{121323} = \frac{(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})^2}{12t^4} = -(W \cdot C)_{122313}, \]

\[ (W \cdot C)_{143413} = \frac{\sin^2 \theta \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi} \right)^2}{12t(t - 2\xi)} = -(W \cdot C)_{133414}, \]

\[ (W \cdot C)_{121424} = \frac{\sin^2 \theta \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi} \right)^2}{12t^4} = -(W \cdot C)_{122414}, \]

\[ (W \cdot C)_{233424} = \frac{(t - 2\xi) \sin^2 \theta \left( t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi} \right)^2}{12t^3} = -(W \cdot C)_{243423}; \]

\[ (W \cdot P)_{121323} = -(W \cdot P)_{123123} = -(W \cdot P)_{132123} = \frac{(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\dot{\xi})}{18t^4}, \]

\[ (W \cdot P)_{133313} = \frac{1}{\sin^4 \theta} (W \cdot P)_{144414} = \frac{(t\ddot{\xi} - 2\dot{\xi})(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})}{18(t - 2\xi)}, \]

\[ (W \cdot P)_{143413} = (W \cdot P)_{134314} = \frac{\sin^2 \theta \left( t(t\ddot{\xi} - 5\dot{\xi}) + 9\dot{\xi} \right)(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})}{18t(t - 2\xi)}, \]

\[ (W \cdot P)_{144313} = (W \cdot P)_{133414} = (W \cdot P)_{344113} = \frac{\sin^2 \theta (t\ddot{\xi} - 3\dot{\xi})(t(t\ddot{\xi} - 4\dot{\xi}) + 6\dot{\xi})}{6t(t - 2\xi)}, \]
\[(W \cdot P)_{122414} = -(W \cdot P)_{121424} = -\frac{\sin^2 \theta \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \ddot{\xi} \right) \left(t(t \ddot{\xi} - 3 \dot{\xi}) + 3 \xi \right)}{6t^4},\]

\[(W \cdot P)_{124214} = -(W \cdot P)_{142124} = -(W \cdot P)_{124124} = -\frac{\sin^2 \theta \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \ddot{\xi} \right) \left(t(2t \ddot{\xi} - 7 \dot{\xi}) + 9 \xi \right)}{18t^4},\]

\[(W \cdot P)_{233323} = \frac{1}{\sin^4 \theta} (W \cdot P)_{244424} = -\frac{(t - 2 \xi)(t \ddot{\xi} - 2 \dot{\xi}) \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \dot{\xi} \right)}{18t^2},\]

\[(W \cdot P)_{243423} = (W \cdot P)_{234423} = (W \cdot P)_{233424} = -\frac{(t - 2 \xi) \sin^2 \theta \left(t(t \ddot{\xi} - 3 \dot{\xi}) \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \dot{\xi} \right) \right)}{6t^3}.\]

For the tensors \(K \cdot R, K \cdot S, K \cdot C, K \cdot W, K \cdot K\) and \(K \cdot P\) we have the following relations:

\[(K \cdot R)_{121323} = \frac{(t^2 \ddot{\xi} + 2 \dot{\xi}) \left(t(t \ddot{\xi} - 3 \dot{\xi}) + 3 \xi \right)}{2t^4} = -(K \cdot R)_{122313},\]

\[(K \cdot R)_{133414} = \frac{\sin^2 \theta(t \ddot{\xi} - 3 \dot{\xi})(t^2 \ddot{\xi} + 2 \dot{\xi})}{2t(t - 2 \xi)} = -(K \cdot R)_{134314},\]

\[(K \cdot R)_{121424} = \frac{\sin^2 \theta(t^2 \ddot{\xi} + 2 \dot{\xi}) \left(t(t \ddot{\xi} - 3 \dot{\xi}) + 3 \xi \right)}{2t^4} = -(K \cdot R)_{122414},\]

\[(K \cdot R)_{243423} = \frac{(t - 2 \xi) \sin^2 \theta(t \ddot{\xi} - 3 \dot{\xi})(t^2 \ddot{\xi} + 2 \dot{\xi})}{2t^3} = -(K \cdot R)_{233424},\]

\[(K \cdot S)_{1313} = \frac{(t^2 \ddot{\xi} + 2 \dot{\xi})(t \ddot{\xi} - 2 \dot{\xi})}{2t^2(t - 2 \xi)},\quad (K \cdot S)_{1414} = \frac{\sin^2 \theta(t^2 \ddot{\xi} + 2 \dot{\xi})(t \ddot{\xi} - 2 \dot{\xi})}{2t^2(t - 2 \xi)},\]

\[(K \cdot S)_{2323} = \frac{(t - 2 \xi)(t^2 \ddot{\xi} + 2 \dot{\xi})(t \ddot{\xi} - 2 \dot{\xi})}{2t^4},\quad (K \cdot S)_{2424} = \frac{(t - 2 \xi) \sin^2 \theta(t^2 \ddot{\xi} + 2 \dot{\xi})(t \ddot{\xi} - 2 \dot{\xi})}{2t^4};\]

\[(K \cdot C)_{121323} = \frac{(t^2 \ddot{\xi} + 2 \dot{\xi}) \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \ddot{\xi} \right)}{4t^4} = -(K \cdot C)_{122313},\]

\[(K \cdot C)_{143413} = \frac{\sin^2 \theta(t^2 \ddot{\xi} + 2 \dot{\xi}) \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \ddot{\xi} \right)}{4t(t - 2 \xi)} = -(K \cdot C)_{134314},\]

\[(K \cdot C)_{121424} = \frac{\sin^2 \theta(t^2 \ddot{\xi} + 2 \dot{\xi}) \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \ddot{\xi} \right)}{4t^4} = -(K \cdot C)_{122414},\]

\[(K \cdot C)_{233424} = \frac{(t - 2 \xi) \sin^2 \theta(t^2 \ddot{\xi} + 2 \dot{\xi}) \left(t(t \ddot{\xi} - 4 \dot{\xi}) + 6 \ddot{\xi} \right)}{4t^3} = -(K \cdot C)_{243423};\]
\[ (K \cdot P)_{12313} = -(K \cdot P)_{121323} = \frac{1}{\sin^2 \theta} (K \cdot P)_{122414} \]
\[ = -\frac{1}{\sin^2 \theta} (K \cdot P)_{1214424} = -\frac{(t^2 \dot{\xi} + 2\xi) \left(t(2t\ddot{\xi} - 3\dot{\xi}) + 3\dot{\xi}\right)}{2t^4}, \]
\[ (K \cdot P)_{123213} = -(K \cdot P)_{132123} = -(K \cdot P)_{123123} = \frac{(t^2 \dot{\xi} + 2\xi) \left(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\dot{\xi}\right)}{6t^4}, \]
\[ (K \cdot P)_{133313} = \frac{1}{\sin^4 \theta} (K \cdot P)_{144414} = \frac{(t^2 \dot{\xi} + 2\xi) \left(t(\ddot{\xi} - 2\dot{\xi})\right)}{6(t - 2\xi)}, \]
\[ (K \cdot P)_{143413} = (K \cdot P)_{134314} = \frac{\sin^2 \theta(t^2 \dot{\xi} + 2\xi) \left(t(t\ddot{\xi} - 5\dot{\xi}) + 9\dot{\xi}\right)}{6t(t - 2\xi)}, \]
\[ (K \cdot P)_{143313} = (K \cdot P)_{133414} = (K \cdot P)_{344113} = \frac{\sin^2 \theta(t\dot{\xi} - 3\xi)(t^2 \dot{\xi} + 2\xi)}{2t(t - 2\xi)}, \]
\[ (K \cdot P)_{233323} = \frac{1}{\sin^4 \theta} (K \cdot P)_{244424} = -\frac{(t - 2\xi)(t^2 \dot{\xi} + 2\xi)(t\ddot{\xi} - 2\dot{\xi})}{6t^2}, \]
\[ (K \cdot P)_{244323} = (K \cdot P)_{344223} = (K \cdot P)_{234244} = -\frac{(t - 2\xi)\sin^2 \theta(t\dot{\xi} - 3\xi)(t^2 \dot{\xi} + 2\xi)}{2t^3}, \]
\[ (K \cdot P)_{124124} = (K \cdot P)_{142124} = -(K \cdot P)_{124124} = -\frac{\sin^2 \theta(t^2 \dot{\xi} + 2\xi) \left(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\dot{\xi}\right)}{6t^4}, \]
\[ (K \cdot P)_{234324} = (K \cdot P)_{243423} = -\frac{(t - 2\xi)\sin^2 \theta(t^2 \dot{\xi} + 2\xi) \left(t(t\ddot{\xi} - 5\dot{\xi}) + 9\dot{\xi}\right)}{6t^3}. \]

For the tensors \( P \cdot R, P \cdot S, P \cdot C, P \cdot W, P \cdot K \) and \( P \cdot P \) we have the following relations:

\[ (P \cdot R)_{121323} = \frac{t^2 \dot{\xi} \left(5\dot{\xi} - 2t\ddot{\xi}\right) + 2t\xi \left(2t\ddot{\xi} - 7\dot{\xi}\right) + 9\dot{\xi}^2}{3t^4} = (P \cdot R)_{122331}, \]
\[ (P \cdot R)_{121424} = \frac{\sin^2 \theta \left(t^2 \dot{\xi} \left(5\dot{\xi} - 2t\ddot{\xi}\right) + 2t\xi \left(2t\ddot{\xi} - 7\dot{\xi}\right) + 9\dot{\xi}^2\right)}{3t^4} = (P \cdot R)_{124214}, \]
\[ (P \cdot R)_{134313} = \frac{\sin^2 \theta \left(t^2 \dot{\xi}^2 + 2t\xi \left(t\ddot{\xi} - 5\dot{\xi}\right) + 9\dot{\xi}^2\right)}{3t(t - 2\xi)} = (P \cdot R)_{134314}, \]
\[ (P \cdot R)_{233424} = \frac{(t - 2\xi)\sin^2 \theta \left(t^2 \dot{\xi}^2 + 2t\xi \left(t\ddot{\xi} - 5\dot{\xi}\right) + 9\dot{\xi}^2\right)}{3t^3} = (P \cdot R)_{244323}, \]
\[ (P \cdot S)_{1313} = \frac{(\xi - t\dot{\xi})(t\ddot{\xi} - 2\dot{\xi})}{t^2(t - 2\xi)} = -(P \cdot S)_{1313}, \]
\[(P \cdot S)_{14114} = \frac{\sin^2 \theta (\xi - t\dot{\xi})(t\ddot{\xi} - 2\dot{\xi})}{t^2(t - 2\xi)} = -(P \cdot S)_{1441},\]
\[(P \cdot S)_{2323} = \frac{(t - 2\xi)(t\ddot{\xi} - \xi)(t\ddot{\xi} - 2\dot{\xi})}{t^4} = -(P \cdot S)_{2332},\]
\[(P \cdot S)_{2424} = \frac{(t - 2\xi)\sin^2 \theta(t\ddot{\xi} - \xi)(t\ddot{\xi} - 2\dot{\xi})}{t^4} = -(P \cdot S)_{2442};\]
\[(P \cdot C)_{121323} = \frac{t(t\ddot{\xi} - 5\ddot{\xi}) + 9\xi}{18t^4} \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right) = (P \cdot C)_{123213},\]
\[(P \cdot C)_{121424} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 5\dot{\xi}) + 9\xi\right)}{18t^4} \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right) = (P \cdot C)_{124214},\]
\[(P \cdot C)_{134341} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{18t(t - 2\xi)} \left(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\xi\right) = (P \cdot C)_{143413},\]
\[(P \cdot C)_{233424} = \frac{(t - 2\xi)\sin^2 \theta \left(t(t\ddot{\xi} - 4\dot{\xi}) + 6\xi\right)}{18t^3} \left(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\xi\right) = (P \cdot C)_{244323};\]
\[(P \cdot P)_{121323} = \frac{t(t\ddot{\xi} - 5\ddot{\xi}) + 9\xi}{9t^4} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = (P \cdot P)_{232113},\]
\[(P \cdot P)_{121424} = \frac{\sin^2 \theta \left(t(t\ddot{\xi} - 5\dot{\xi}) + 9\xi\right)}{9t^4} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = (P \cdot P)_{242114},\]
\[(P \cdot P)_{123123} = \frac{(t\ddot{\xi} - 3\ddot{\xi})}{3t^4} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = (P \cdot P)_{231213},\]
\[(P \cdot P)_{124124} = \frac{\sin^2 \theta \left(t\ddot{\xi} - 3\ddot{\xi}\right)}{3t^4} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = (P \cdot P)_{241214},\]
\[(P \cdot P)_{131113} = \frac{-(t\ddot{\xi} - 2\ddot{\xi})}{9t(t - 2\xi)^2} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = -(P \cdot P)_{131131},\]
\[(P \cdot P)_{133331} = \frac{9(t\ddot{\xi} - 3\ddot{\xi})}{9(t - 2\xi)} \left(t\ddot{\xi} - 2\dot{\xi}\right) = -(P \cdot P)_{133331},\]
\[(P \cdot P)_{133441} = \frac{\sin^2 \theta \left(t\ddot{\xi} - 3\ddot{\xi}\right)}{9t(t - 2\xi)} \left(t(2t\ddot{\xi} - 7\dot{\xi}) + 9\xi\right) = (P \cdot P)_{344113},\]
\[(P \cdot P)_{134341} = \frac{\sin^2 \theta \left(t\ddot{\xi} - 3\ddot{\xi}\right)}{3t(t - 2\xi)} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = (P \cdot P)_{341431},\]
\[(P \cdot P)_{141141} = \frac{\sin^2 \theta \left(t\ddot{\xi} - 2\ddot{\xi}\right)}{9t(t - 2\xi)^2} \left(t(t\ddot{\xi} - 3\dot{\xi}) + 3\xi\right) = -(P \cdot P)_{141114}.\]
For the tensors $Q, S, C, P$ we have the following relations:

\[ (P \cdot P)_{144441} = \frac{\sin^4 \theta \left( t \dot{\xi} - 3 \dot{\xi} \right) \left( t \ddot{\xi} - 2 \ddot{\xi} \right)}{9(t - 2\xi)} = -(P \cdot P)_{144441}, \]

\[ (P \cdot P)_{232232} = \frac{(t - 2\xi)^2 \left( t \ddot{\xi} - 2 \ddot{\xi} \right) \left( t \dot{\xi} - 3 \dot{\xi} \right)}{9t^5} = -(P \cdot P)_{232232}, \]

\[ (P \cdot P)_{233332} = \frac{(t - 2\xi) \left( t \ddot{\xi} - 3 \ddot{\xi} \right) \left( t \dot{\xi} - 2 \dot{\xi} \right)}{9t^2} = -(P \cdot P)_{233332}, \]

\[ (P \cdot P)_{234424} = \frac{(t - 2\xi) \sin^2 \theta \left( t \ddot{\xi} - 3 \ddot{\xi} \right) \left( t \dot{\xi} - 3 \dot{\xi} \right)}{9t^3} = (P \cdot P)_{234424}, \]

\[ (P \cdot P)_{234324} = \frac{(t - 2\xi) \sin^2 \theta \left( t \ddot{\xi} - 3 \ddot{\xi} \right) \left( t \dot{\xi} - 3 \dot{\xi} \right)}{3t^3} = (P \cdot P)_{234324}, \]

\[ (P \cdot P)_{242224} = \frac{(t - 2\xi)^2 \sin^2 \theta \left( t \ddot{\xi} - 2 \ddot{\xi} \right) \left( t \dot{\xi} - 3 \dot{\xi} \right)}{9t^5} = -(P \cdot P)_{242224}, \]

\[ (P \cdot P)_{244442} = \frac{(t - 2\xi) \sin^4 \theta \left( t \ddot{\xi} - 3 \ddot{\xi} \right) \left( t \ddot{\xi} - 2 \ddot{\xi} \right)}{9t^2} = -(P \cdot P)_{244442}. \]
\[ Q(S,W)_{121332} = \frac{t^3 \dddot{\xi}^2 + 4\dot{\xi} \left( 6\xi - 7t\dot{\xi} \right) + 6t \left( t\dddot{\xi} + \dot{\xi} \right) \dddot{\xi}}{6t^3} = Q(S,W)_{122313}, \]
\[ Q(S,W)_{121442} = \frac{\sin^2 \theta \left( t^3 \dddot{\xi}^2 + 4\dot{\xi} \left( 6\xi - 7t\dot{\xi} \right) + 6t \left( t\dddot{\xi} + \dot{\xi} \right) \dddot{\xi} \right)}{6t^3} = Q(S,W)_{122414}, \]
\[ Q(S,W)_{133414} = \frac{\sin^2 \theta \left( 12\xi \left( t\dddot{\xi} + \dot{\xi} \right) - t \left( t^2 \dddot{\xi}^2 + 8\dot{\xi}^2 \right) \right)}{6(t - 2\xi)} = Q(S,W)_{143431}, \]
\[ Q(S,W)_{133441} = \frac{\sin^2 \theta \left( t^3 \dddot{\xi}^2 + 8t\dot{\xi}^2 - 12\xi \left( t\dddot{\xi} + \dot{\xi} \right) \right)}{6(t - 2\xi)} = Q(S,W)_{143413}, \]
\[ Q(S,W)_{233424} = \frac{(t - 2\xi) \sin^2 \theta \left( t^3 \dddot{\xi}^2 + 8t\dot{\xi}^2 - 12\xi \left( t\dddot{\xi} + \dot{\xi} \right) \right)}{6t^2} = Q(S,W)_{243432}, \]
\[ Q(S,K)_{121332} = \frac{t^3 \dddot{\xi}^2 + 2t\xi \dddot{\xi} + 8\dot{\xi} \left( \xi - t\dot{\xi} \right)}{2t^3} = Q(S,K)_{122313}, \]
\[ Q(S,K)_{121442} = \frac{\sin^2 \theta \left( t^3 \dddot{\xi}^2 + 2t\xi \dddot{\xi} + 8\dot{\xi} \left( \xi - t\dot{\xi} \right) \right)}{2t^3} = -Q(S,K)_{242114}, \]
\[ Q(S,K)_{133414} = \frac{\sin^2 \theta \left( 2\xi \dddot{\xi} + t \left( 2\xi - t\dddot{\xi} \right) \dot{\xi} \right)}{t - 2\xi} = Q(S,K)_{341431}, \]
\[ Q(S,K)_{133441} = \frac{\sin^2 \theta \left( t \left( t\dddot{\xi} - 2\xi \right) \dddot{\xi} - 2\xi \dddot{\xi} \right)}{t - 2\xi} = Q(S,K)_{344131}, \]
\[ Q(S,K)_{233424} = \frac{(t - 2\xi) \sin^2 \theta \left( t \left( t\dddot{\xi} - 2\xi \right) \dddot{\xi} - 2\xi \dddot{\xi} \right)}{t^2} = Q(S,K)_{342432}, \]
\[ Q(S,P)_{121332} = \left( t\dddot{\xi} + 4\dot{\xi} \right) \left( t \left( t\dddot{\xi} - 3\dddot{\xi} \right) + 3\xi \right) = Q(S,P)_{122313} = -Q(S,P)_{131223}, \]
\[ Q(S,P)_{121442} = \frac{\sin^2 \theta \left( t\dddot{\xi} + 4\dot{\xi} \right) \left( t \left( t\dddot{\xi} - 3\dddot{\xi} \right) + 3\xi \right)}{3t^3} = Q(S,P)_{122414} = -Q(S,P)_{141224}, \]
\[ Q(S,P)_{123123} = \frac{4\dot{\xi} \left( \xi - t\dot{\xi} \right) + t \left( t\dddot{\xi} + \dot{\xi} \right) \dddot{\xi}}{t^3} = Q(S,P)_{132132} = Q(S,P)_{231231}, \]
\[ Q(S,P)_{124124} = \frac{\sin^2 \theta \left( 4\dot{\xi} \left( \xi - t\dot{\xi} \right) + t \left( t\dddot{\xi} + \dot{\xi} \right) \dddot{\xi} \right)}{t^3} = Q(S,P)_{124214} = Q(S,P)_{142142}, \]
\[ Q(S,P)_{131113} = \frac{t\dddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3(t - 2\xi)^2} = -Q(S,P)_{131131}, \]
\[ Q(S, P)_{13331} = \frac{2t\ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3(t - 2\xi)} = -Q(S, P)_{13313}, \]

\[ Q(S, P)_{133441} = \frac{2\sin^2 \theta \left( t\ddot{\xi} - 3\xi \right) \left( t\dddot{\xi} + \dot{\xi} \right)}{3(t - 2\xi)} = -Q(S, P)_{144313} = Q(S, P)_{343114}, \]

\[ Q(S, P)_{144441} = \frac{2\sin^2 \theta \left( \xi \left( t\dddot{\xi} + \dot{\xi} \right) - t\ddot{\xi}^2 \right)}{t - 2\xi} = -Q(S, P)_{134314} = -Q(S, P)_{341413}, \]

\[ Q(S, P)_{141114} = \frac{t\sin^2 \theta \ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3(t - 2\xi)^2} = -Q(S, P)_{141141}, \]

\[ Q(S, P)_{144444} = \frac{2t\sin^4 \theta \ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3(t - 2\xi)} = -Q(S, P)_{144414}, \]

\[ Q(S, P)_{232223} = \frac{(t - 2\xi)^2 \ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3t^3} = -Q(S, P)_{232232}, \]

\[ Q(S, P)_{233332} = \frac{2(t - 2\xi)\ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3t} = -Q(S, P)_{233332}, \]

\[ Q(S, P)_{234324} = \frac{2(t - 2\xi) \sin^2 \theta \left( t\ddot{\xi} - 3\xi \right) \left( t\dddot{\xi} + \dot{\xi} \right)}{3t^2} = Q(S, P)_{343242} = -Q(S, P)_{344232}, \]

\[ Q(S, P)_{234324} = \frac{2(t - 2\xi) \sin^2 \theta \left( \xi \left( t\dddot{\xi} + \dot{\xi} \right) - t\ddot{\xi}^2 \right)}{t^2} = -Q(S, P)_{342324} = -Q(S, P)_{243432}, \]

\[ Q(S, P)_{242224} = \frac{(t - 2\xi)^2 \sin^2 \theta \ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3t^3} = -Q(S, P)_{242242}, \]

\[ Q(S, P)_{244424} = \frac{2(t - 2\xi) \sin^4 \theta \ddot{\xi} \left( t\dddot{\xi} - 2\dot{\xi} \right)}{3t} = -Q(S, P)_{244442}. \]

6. Conclusion

From the above results and discussion we conclude that Deszcz symmetric spaces are geometric models of the interior black hole spacetime and hence the defining conditions of Deszcz symmetric spaces are, physically, very stronger. However, for a specific value of \( \xi \) the interior black hole spacetime turns into a semisymmetric spacetime and also a generalized Ricci pseudosymmetric spacetime.
ACKNOWLEDGEMENT

The work was carried out when the first named author visited Department of Mathematics of the Sar- dar Patel University as a visiting fellow under their UGC-SAP-DRS programme (F-510/5/DRS/2009 (SAP-II)). He also greatfully acknowledges the financial support of CSIR, New Delhi, India [Project F. No. 25(0171)/09/EMR-II]. The second named author is supported by a grant of the Technische Universität Berlin (Germany).

REFERENCES

1. A. Adamów and R. Deszcz, On totally umbilical submanifolds of some class of Riemannian manifolds, *Demonstr. Math.*, **16** (1983), 39-59.

2. A. C. Asperti, G. A. Lobos, and F. Mercuri, Pseudo-parallel submanifolds of a space form, *Adv. Geom.*, **2** (2002), 57-71.

3. M. Berger and D. Ebin, Some characterizations of the space of symmetric tensors on a Riemannian manifold, *J. Diff. Geom.*, **3** (1969), 379-392.

4. A. L. Besse, Einstein manifolds, *Ergeb. Math. Grenzgeb.*, 3. Folge. Bd. 10, Springer-Verlag, Berlin, Heidelberg, New York, 1987.

5. R. L. Bishop and B. O’Neill, Manifolds of negative curvature, *Trans. Amer. Math. Soc.*, **145** (1969), 1-49.

6. M. Blau, *Lecture notes on general relativity*, http://www.blau.itp.unibe.ch/Lecturenotes.html, 2018.

7. J. P. Bourguignon, Les variétés de dimension 4 à signature non nulle dont la courbure est harmonique sont d’Einstein, *Invent. Math.*, **63**(2) (1981), 263-286.

8. M. Cahen and M. Parker, Sur des classes d’espaces pseudo-riemanniens symetriques, *Bull. Soc. Math. Belg.*, **22** (1970), 339-354.

9. M. Cahen and M. Parker, Pseudo-Riemannian symmetric spaces, *Mem. Amer. Math. Soc.*, **24** (1980), 1-108.

10. B. J. Carr, *Black hole in cosmology and astrophysics*, Proc. of the 46 Scottish Univ. Summer School in Phys., Aberdeen, July 1995, NATO Advanced Study Institute and IOP, London, 143-202.

11. Cartan, É., Sur une classe remarquable désespaces de Riemann, I, *Bull. de la Soc. Math. de France*, **54** (1926), 214-216.

12. Cartan, É., Sur une classe remarquable désespaces de Riemann, II, *Bull. de la Soc. Math. de France*, **55** (1927), 114-134.

13. Cartan, É., *Leçons sur la géométrie des espaces de Riemann*, Gauthier Villars, 2nd ed., Paris, 1946.
14. M. C. Chaki, On pseudosymmetric manifolds, *An. Ştiinţ. ale Univ., AL I. Cuza din Iaşi N. Ser. Sect. Ia*, **33** (1987), 53-58.

15. B.-Y. Chen, *Differential Geometry of Warped Product Manifolds and Submanifolds*, World Sci., 2017.

16. J. Chojnacka-Dulas, R. Deszcz, M. Głogowska, and M. Prvanović, On warped product manifolds satisfying some curvature conditions, *J. Geom. Phys.*, **74** (2013), 328-341.

17. S. Decu, M. Petrović-Torgašev, A. Šebeković, and L. Verstraelen, On the Roter type of Wintgen ideal submanifolds, *Rev. Roumaine Math. Pures Appl.*, **57** (2012), 75-90.

18. F. Defever, R. Deszcz, M. Hotloš, M. Kucharski, and Z. Şentürk, Generalisations of Robertson-Walker spaces, *Ann. Univ. Sci. Budapest. Eötvös Sect. Math.*, **43** (2000), 13-24.

19. F. Defever, R. Deszcz, Z. Şentürk, L. Verstraelen, and S. Yaprak, P. J. Ryan’s problem in semi-Riemannian space forms, *Glasgow Math. J.*, **41** (1999), 271-281.

20. F. Defever, R. Deszcz, L. Verstraelen, and L. Vrancken, On pseudosymmetric space-times, *J. Math. Phys.*, **35** (1994), 5908-5921.

21. J. Deprez, W. Roter, and L. Verstraelen, Conditions on the projective curvature tensor of conformally flat Riemannian manifolds, *Kyungpook Math. J.*, **29** (1989), 153-165.

22. A. Derdziński, Some remarks on the local structure of Codazzi tensors, *Glob. Diff. Geom. Glob. Ann.*, lecture notes, Springer-Verlag, **838** (1981), 251-255.

23. A. Derdziński, On compact Riemannian manifolds with harmonic curvature, *Math. Ann.*, **259** (1982), 145-152.

24. A. Derdziński and C. L. Shen, Codazzi tensor fields, curvature and Pontryagin forms, *Proc. London Math. Soc.*, **47**(3) (1983), 15-26.

25. R. Deszcz, Notes on totally umbilical submanifolds, Geometry and Topology of Submanifolds, Luminy, May 1987, *World Sci. Publ.*, Singapore (1989), 89-97.

26. R. Deszcz, On four-dimensional warped product manifolds satisfying certain pseudosymmetry curvature conditions, *Colloq. Math.*, **62** (1991), 103-120.

27. R. Deszcz, On pseudosymmetric totally umbilical submanifolds of Riemannian manifolds admitting some types of generalized curvature tensors, *Zesz. Nauk. Politechn. Slask.*, Issue dedicated to the 70th Birthday of Prof. Mieczysław Kucharzewski, **68** (1993), 171-187.

28. R. Deszcz, On pseudosymmetric spaces, *Bull. Belg. Math. Soc., Ser. A*, **44** (1992), 1-34.

29. R. Deszcz, F. Dillen, L. Verstraelen, and L. Vrancken, Quasi-Einstein totally real submanifolds of the nearly Kähler 6-sphere, *Tôhoku Math. J.*, **51** (1999), 461-478.

30. R. Deszcz and M. Głogowska, Some examples of nonsemisymmetric Ricci-semisymmetric hypersurfaces, *Colloq. Math.*, **94** (2002), 87-101.
31. R. Deszcz, M. Głogowska, H. Hashiguchi, M. Hotloś, and M. Yawata, On semi-Riemannian manifolds satisfying some conformally invariant curvature condition, *Colloq. Math.*, **131** (2013), 149-170.

32. R. Deszcz, M. Głogowska, and M. Hotloś, Some identities on hypersurfaces in conformally flat spaces, in: *Proceedings of the International Conference XVI Geometrical Seminar*, September, 20-25, Vrnjačka banja, September, 20-25, 2010, Faculty of Science and Mathematics, University of Niš, Serbia, 2011, 34-39.

33. R. Deszcz, M. Głogowska, M. Hotloś, and M. Sawicz, A Survey on Generalized Einstein Metric Conditions, *Advances in Lorentzian Geometry: Proceedings of the Lorentzian Geometry Conference in Berlin*, AMS/IP Studies in Advanced Mathematics **49**, S.-T. Yau (series ed.), M. Plaue, A.D. Rendall and M. Scherfner (eds.) (2011), 27-46.

34. R. Deszcz, M. Głogowska, M. Hotloś, and Z. Sentürk, On certain quasi-Einstein semi-symmetric hypersurfaces, *Ann. Univ. Sci. Budapest. Eötvös Sect. Math.*, **41** (1998), 151-164.

35. R. Deszcz, M. Głogowska, M. Hotloś, and G. Zafindratafa, Hypersurfaces in spaces forms satisfying some curvature conditions, *J. Geom. Phys.*, **99** (2016), 218-231.

36. R. Deszcz, M. Głogowska, J. Jełowicki, and G. Zafindratafa, Curvature properties of some class of warped product manifolds, *Int. J. Geom. Meth. Modern Phys.*, **13** (2016), 1550135 (36 pages).

37. R. Deszcz, M. Głogowska, M. Petrović-Torgašev, and L. Verstraelen, Curvature properties of some class of minimal hypersurfaces in Euclidean spaces, *Filomat*, **29** (2015), 479-492.

38. R. Deszcz and W. Grycak, On some class of warped product manifolds, *Bull. Inst. Math. Acad. Sinica*, **15** (1987), 311-322.

39. R. Deszcz, S. Haesen, and L. Verstraelen, Classification of space-times satisfying some pseudo-symmetry type conditions, *Soochow J. Math.*, **23** (2004), 339-349 (Special issue in honor of Professor Bang-Yen Chen).

40. R. Deszcz, S. Haesen, and L. Verstraelen, On natural symmetries, *Topics in Differential Geometry*, Eds. A. Mihai, I. Mihai and R. Miron, Editura Academiei Române, 2008.

41. R. Deszcz M. and Hotloś, Remarks on Riemannian manifolds satisfying certain curvature condition imposed on the Ricci tensor, *Pr. Nauk. Politech. Szczec.*, **11** (1989), 23-34.

42. R. Deszcz, M. and Hotloś, On some pseudosymmetry type curvature condition, *Tsukuba J. Math.*, **27** (2003), 13-30.

43. R. Deszcz and M. Hotloś, On hypersurfaces with type number two in spaces of constant curvature, *Ann. Univ. Sci. Budapest. Eötvös Sect. Math.*, **46** (2003), 19-34.

44. R. Deszcz and M. Hotloś, On geodesic mappings in particular class of Roter spaces, arXiv: 1812.00670 [math.DG], submitted 3 December 2018.

45. R. Deszcz, M. Hotloś, J. Jełowicki, H. Kundu, and A. A. Shaikh, Curvature properties of Gödel metric,
46. R. Deszcz, M. Hotloś, J. Jelowicki, H. Kundu, and A. A. Shaikh, Erratum - Curvature properties of G"odel metric, *Int. J. Geom. Methods Mod. Phys.*, **16** (2019), 1992002 (4 pages).

47. R. Deszcz, M. Hotloś, and Z. Şentürk, Quasi-Einstein hypersurfaces in semi-Riemannian space forms, *Colloq. Math.*, **81** (2001), 81-97.

48. R. Deszcz, M. Hotloś, and Z. Şentürk, On curvature properties of quasi-Einstein hypersurfaces in semi-Euclidean spaces, *Soochow J. Math.*, **27** (2001), 375-389.

49. R. Deszcz and D. Kowalczyk, On some class of pseudosymmetric warped products, *Colloq. Math.*, **97** (2003), 7-22.

50. R. Deszcz, M. Plaue, and M. Scherfner, On Roter type warped products with 1-dimensional fibers, *J. Geom. Phys.*, **69** (2013), 1-11.

51. R. Deszcz and M. Scherfner, On a particular class of warped products with fibres locally isometric to generalized Cartan hypersurfaces, *Colloq. Math.*, **109** (2007), 13-29.

52. R. Deszcz, P. Verheyen, and L. Verstraelen, On some generalized Einstein metric conditions, *Publ. Inst. Math. (Beograd) (N.S.)*, **60** (**74**) (1996), 108-120.

53. R. Deszcz, L. Verstraelen, and L. Vrancken, The symmetry of warped product space-times, *Gen. Rel. Gravitation*, **23** (1991), 671-681.

54. R. Deszcz, L. Verstraelen, and S. Yaprak, Warped products realizing a certain condition of pseudosymmetry type imposed on the Weyl curvature tensor, *Chinese J. Math.*, **22** (1994), 139-157.

55. R. Deszcz and S. Yaprak, Curvature properties of certain pseudosymmetric manifolds, *Publ. Math. Debrecen*, **45** (1994), 333-345.

56. R. Doran, F. S. N. Lobo, and P. Crawford, Interior of a Schwarzschild black hole revisited, *Found. Phys.*, **38** (2008), 160-187.

57. D. Dumitru, On quasi-Einstein warped products, *Jordan J. Math. Stat.*, **5** (2012), 85-95.

58. D. Dumitru, A characterization of generalized quasi-Einstein manifolds, *Novi. Sad. J. Math.*, **42** (2012), 89-94.

59. D. Ferus, A remark on Codazzi tensors on constant curvature space, *Glob. Diff. Geom. Glob. Ann.*, Lecture notes, Springer, **838** (1981).

60. M. Głogowska, Semi-Riemannian manifolds whose Weyl tensor is a Kulkarni-Nomizu square, *Publ. Inst. Math. (Beograd) (N.S.)*, **72** (**86**) (2002), 95-106.

61. M. Głogowska, Curvature conditions on hypersurfaces with two distinct principal curvatures, Banach Center Publications, *Inst. Math. Polish Acad. Sci.*, **69** (2005), 133-143.

62. M. Głogowska, On quasi-Einstein Cartan type hypersurfaces, *J. Geom. Phys.*, **58** (2008), 599-614.
63. M. Głogowska, On Roter type manifolds, *Pure and Applied Differential Geometry - PADGE 2007*, Berichte aus der Mathematik, Shaker Verlag, Aachen (2007), 114-122.

64. A. Gray, Einstein-like manifolds which are not Einstein, *Geom. Dedicata*, 7 (1978), 259-280.

65. J. B. Griffiths and J. Podolský, *Exact space-times in Einsteins general relativity*, Cambridge Univ. Press, 2009.

66. S. Haesen and L. Verstraelen, Classification of the pseudosymmetric space-times, *J. Math. Phys.*, 45 (2004), 2343-2346.

67. S. Haesen and L. Verstraelen, Properties of a scalar curvature invariant depending on two planes, *Manuscripta Math.*, 122 (2007), 59-72.

68. S. Haesen and L. Verstraelen, Natural intrinsic geometrical symmetries, *Symmetry, Integrability and Geometry: Methods and Applications (SIGMA)*, 5 (2009), 086, 14 pages.

69. A. H. Hasmani and V. G. Khambholja, Algebraic computations of Riemannian curvature tensor for 5D space using Mathematica, National Conference on Recent Trends in Engineering and Technology, BVM Engineering College, 13-14 May, 2011.

70. A. H. Hasmani and V. G. Khambholja, Interior black-hole solution with anisotropic fluid, *J. Dyna. Sys. Geom. Th.*, 10(1) (2012), 47-52.

71. A. H. Hasmani and V. G. Khambholja, A metric for 5D interior black-hole solution, *J. Dyna. Sys. Geom. Th.*, 10(2) (2012), 125-128.

72. B. Jahanara, S. Haesen, Z. Sentürk, and L. Verstraelen, On the parallel transport of the Ricci curvatures, *J. Geom. Phys.*, 57 (2007), 1771-1777.

73. D. Kowalczyk, On some subclass of semi-symmetric manifolds, *Soochow J. Math.*, 27 (2001), 445-461.

74. D. Kowalczyk, On the Reissner-Nordström-de Sitter type spacetimes, *Tsukuba J. Math.*, 30 (2006), 363-381.

75. O. Kowalski and M. Sekigawa, Pseudo-symmetric spaces of constant type in dimension three - elliptic spaces, *Rend. Mat. Appl.*, 7 (17) (1997), 477-512.

76. O. Kowalski and M. Sekigawa, Pseudo-symmetric spaces of constant type in dimension three - non-elliptic spaces, *Bull. Tokyo Gakugei Univ., Sect. IV, Math. Nat. Sci.*, 50 (1998), 1-28.

77. G. I. Kruchkovich, On some class of Riemannian spaces, *Trudy Sem. Vektor. Tensor. Anal.*, 11 (1961) 103-128 (in Russian).

78. G. I. Kruchkovich, *On spaces of V.F. Kagan*, in: Subprojective Spaces, ed. V.F. Kagan, (Fiznatgiz, Moscow, 1961), pp. 163-198 (in Russian).

79. A. Lichnerowicz, Courbure, nombres de Betti, et espaces symetriques, *Proc. Int. Cong. Math.*, 2 (1952), 216-223.
80. B. O’ Neill, *Semi-Riemannian Geometry with application to the relativity*, Academic Press, New York, London, 1983.
81. J. R. Oppenheimer and H. Snyder, On Continued Gravitational Contraction, *Phys. Rev.*, 56 (1939), 455-459.
82. R. Penrose, Gravitational Collapse: The Role of General Relativity, *Riv. Nuovo Vimento Soc. Ital. Fis.*, 1 (1969), 252-276.
83. V. Perlick, Gravitational Lensing from a Spacetime Perspective, *Living Rev. Relativity*, 7:9 (2004), 1-117. doi: 10.12942/lrr-2004-9.
84. A. Selberg, Harmonic analysis and discontinuous groups in weakly symmetric Riemannian spaces with applications to Dirichlet series, *Indian J. Math. Soc.*, 20 (1956), 47-87.
85. A. A. Shaikh, F. R. Al-Solamy, and I. Roy, On the existence of a new class of semi-Riemannian manifolds, *Math. Sciences*, 7(46) (2013), 1-13.
86. A. A. Shaikh and T. Q. Binh On some class of Riemannian manifolds, *Bull. Transilvania Univ.*, 15(50) (2008), 351-362.
87. A. A. Shaikh, M. Ali, and Z. Ahsan, Curvature properties of Robinson-Trautman metric, *J. of Geom.*, 109(38) (2018), 1-20, doi: 10.1007/s00022-018-0443-1.
88. A. A. Shaikh, T. Q. Binh, and H. Kundu, Curvature properties of generalized pp-wave metric, *Kragujevac J. Math.*, 45(2) (2021), 237-258.
89. A. A. Shaikh, R. Deszcz, M. Hotloš, J. Jełowicki, and H. Kundu, On pseudosymmetric manifolds, *Publ. Math. Debrecen*, 86(3-4) (2015), 433-456.
90. A. A. Shaikh, S. K. Hui, and W. Yoon, W. Dae, On quasi Einstein Spacetimes, *Tsukuba J. Math.*, 33(2) (2009), 305-326.
91. A. A. Shaikh and S. K. Hui, On decomposable quasi-Einstein spaces, *Math. Reports*, 13(63) (2011), 89-94.
92. A. A. Shaikh, Y. H. Kim, and S. K. Hui, On Lorentzian quasi-Einstein manifolds, *J. Korean Math. Soc.*, 48 (2011), 669-689, and Erratum to: On Lorentzian quasi-Einstein manifolds, *J. Korean Math. Soc.*, 48 (2011), 1327-1328.
93. A. A. Shaikh and H. Kundu, On weakly symmetric and weakly Ricci symmetric warped product manifolds, *Publ. Math. Debrecen*, 48(3-4) (2012), 487-505.
94. A. A. Shaikh and H. Kundu, On equivalency of various geometric structures, *J. Geom.*, 105 (2014), 139-165.
95. A. A. Shaikh and H. Kundu, On generalized Roter type manifolds, *Kragujevac J. Math.*, 43(3) (2019), 471-493.
96. A. A. Shaikh and H. Kundu, On warped product generalized Roter type manifolds, *Balkan J. Geom. Appl.*, 21(2) (2016), 82-95.

97. A. A. Shaikh and H. Kundu, On curvature properties of Som-Raychaudhuri spacetime, *J. Geom.*, 108(2) (2016), 501-515.

98. A. A. Shaikh and H. Kundu, On warped product manifolds satisfying some pseudosymmetric type conditions, *Diff. Geom. - Dyn. Syst.*, 19 (2017), 119-135.

99. A. A. Shaikh and H. Kundu, On some curvature restricted geometric structures for projective curvature tensor, *Int. J. Geom. Methods Mod. Phys.*, 15 (2018), 1850157 (38 pages).

100. A. A. Shaikh, H. Kundu, and S. Ali, Md., On warped product super generalized recurrent manifolds, *An. Stiint. Univ. Al. I. Cuza Iasi Mat. (N. S.),* LXIV(1) (2018), 85-99.

101. A. A. Shaikh, H. Kundu, M. Ali, and Z. Ahsan, Curvature properties of a special type of pure radiation metrics, *J. Geom. Phys.*, 136 (2019), 195-206.

102. A. A. Shaikh, H. Kundu, and J. Sen, Curvature properties of Vaidya metric, *Indian J. Math.*, 61(1) (2019), 41-59.

103. A. A. Shaikh and A. Patra, On a generalized class of recurrent manifolds, *Archivum Math.*, 46 (2010), 39-46.

104. A. A. Shaikh and I. Roy, On quasi generalized recurrent manifolds, *Math. Pannonica*, 21 (2010), 251-263.

105. A. A. Shaikh and I. Roy, On weakly generalized recurrent manifolds, *Ann. Univ. Sci. Budapest. Eötvös Sect. Math.*, 54 (2011), 35-45.

106. A. A. Shaikh, I. Roy, and H. Kundu, On some generalized recurrent manifolds, *Bull. Iranian Math. Soc.*, 43(5) (2017), 1209-1225.

107. A. A. Shaikh, I. Roy, and H. Kundu, On the existence of a generalized class of recurrent manifolds, *An. Stiint. Univ. Al. I. Cuza Iasi Mat. (N. S.),* LXIV(2) (2018), 233-251.

108. A. A. Shaikh, D. W. Yoon, and S. K. Hui, On quasi-Einstein spacetimes, *Tsukuba J. Math.*, 33 (2009), 305-326.

109. Y. Simon, Codazzi tensors, *Glob. Diff. Geom. and Glob. Ann.*, lecture notes, Springer-Verlag, 838 (1981), 289-296.

110. Z. I. Szabó, Structure theorems on Riemannian spaces satisfying $R(X, Y) \cdot R = 0$, I, The local version, *J. Diff. Geom.*, 17 (1982), 531-582.

111. Z. I. Szabó, Classification and construction of complete hypersurfaces satisfying $R(X, Y) \cdot R = 0$, *Acta Sci. Math.*, 47 (1984), 321-348.

112. Z. I. Szabó, Structure theorems on Riemannian spaces satisfying $R(X, Y) \cdot R = 0$, II, The global version,
113. L. Támassy, and T. Q. Binh, On weakly symmetric and weakly projective symmetric Riemannian manifolds, *Colloq. Math. Soc. János Bolyai*, 50 (1989), 663-670.

114. L. Verstraelen, *Comments on the pseudo-symmetry in the sense of Ryszard Deszcz*, in: Geometry and Topology of Submanifolds, VI, World Sci., Singapore, 1994, 119-209.

115. L. Verstraelen, A concise mini history of Geometry, *Kragujevac J. Math.*, 38 (2014), 5-21.

116. L. Verstraelen, *Natural extrinsic geometrical symmetries - an introduction*, in: Recent Advances in the Geometry of Submanifolds Dedicated to the Memory of Franki Dillen (1963-2013), Contemporary Mathematics, 674 (2016), 5-16.

117. L. Verstraelen, Foreword, in: B.-Y. Chen, *Differential Geometry of Warped Product Manifolds and Submanifolds*, World Sci., 2017, vii-xxi.

118. A. G. Walker, On Ruse’s spaces of recurrent curvature, *Proc. London Math. Soc.*, 52 (1950), 36-64.