A Global Existence Result for a Semilinear Wave Equation with Lower Order Terms on Compact Lie Groups

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Abstract
In this paper, we study the semilinear wave equation with lower order terms (damping and mass) and with power type nonlinearity $|u|^p$ on compact Lie groups. We will prove the global in time existence of small data solutions in the evolution energy space without requiring any lower bounds for $p > 1$. In our approach, we employ some results from Fourier analysis on compact Lie groups.

Keywords Wave equation · Compact Lie group · Global existence · Group Fourier transform

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1 Introduction
Let $\mathbb{G}$ be a compact, connected Lie group and let $\mathcal{L}$ be the Laplace–Beltrami operator on $\mathbb{G}$ (which coincides with the Casimir element of the enveloping algebra). In the present work, we prove the global existence of small data solutions for the Cauchy problem for the semilinear wave equation with damping and mass and with power type nonlinearity, namely,

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\[
\begin{align*}
\partial_t^2 u - \mathcal{L} u + b \partial_t u + m^2 u &= |u|^p, & x \in \mathbb{G}, & t > 0, \\
u(0, x) &= u_0(x), & x \in \mathbb{G}, \\
\partial_t u(0, x) &= u_1(x), & x \in \mathbb{G},
\end{align*}
\]

(1)

where \( p > 1 \) and \( b, m^2 \) are positive constants.

In [25, 26] the semilinear wave equation in compact Lie groups \((b = m^2 = 0)\) and the semilinear damped wave equation in compact Lie groups \((b > 0, m^2 = 0)\) have been considered, respectively. For both these cases, the nonexistence of global in time solutions can be proved for any \( p > 1 \), under suitable sign assumptions on the Cauchy data. Purpose of this paper is to show that the combined presence of both a damping term \( b \partial_t u \) and of a mass term \( m^2 u \) modifies completely the situation. Indeed, we will show that, as in the Euclidean case (see [27, Chapter 4] and [8]) or, more generally, as for a graded Lie group (cf. [28]) or when one replaces the standard Laplacian with the Dunkl Laplacian (cf. [35]), the solution to the linear homogeneous Cauchy problem

\[
\begin{align*}
\partial_t^2 u - \mathcal{L} u + b \partial_t u + m^2 u &= 0, & x \in \mathbb{G}, & t > 0, \\
u(0, x) &= u_0(x), & x \in \mathbb{G}, \\
\partial_t u(0, x) &= u_1(x), & x \in \mathbb{G},
\end{align*}
\]

(2)

and its derivatives fulfill \( L^2(\mathbb{G}) \)–\( L^2(\mathbb{G}) \) decay estimates with exponential decay rates with respect to the time variable. Using these decay estimates on \( L^2(\mathbb{G}) \) basis, we will be able to prove the existence of a uniquely determined globally in time defined mild solution to (1) provided that \((u_0, u_1)\) has sufficiently small norm in the energy space via a standard contraction argument. Furthermore, a Gagliardo–Nirenberg type inequality recently proved in [30] will be employed to estimate the power nonlinearity in \( L^2(\mathbb{G}) \).

Let us stress that the decay estimates for the linear homogeneous Cauchy problem (2) will be proved via Fourier Analysis on compact Lie groups. In particular, the group Fourier transform and Plancherel formula with respect to the space variables will play a crucial role.

Let us review briefly some results known in literature concerning nonlinear evolution models in compact manifolds. Let us begin by recalling the results obtained by Bourgain in [1, 2] for the nonlinear Schrödinger equation on the torus. In particular, differently from what happens for the corresponding flat case (see [13, 14, 21, 33] and references therein), the Strichartz estimates are not global but only local (in time) in the periodic case. Therefore, in order to get the global in time existence of large data solutions one has to combine a local (in time) existence result with the conservation laws associated with the equation (provided that the nonlinear term satisfies suitable conditions). A further study of the solutions to the nonlinear Schrödinger equation for different (and sometimes more general) compact manifolds is carried in the series of papers by Burq et al. [3–6].

On the other hand, Kapitanski considered in [20] a semilinear wave equation with analogous lower order terms as in (1) in a general compact manifold, but with a quite different class of nonlinear terms (to which the defocusing nonlinearity belongs, for example). In particular, the global existence of large data solutions is obtained
by combining the dissipative effects of the lower order terms with some Strichartz estimates. Nevertheless, the power nonlinearity $|u|^p$ is not in the class of nonlinear terms allowed in [20]. We emphasize that, due to the nonlinear term in (1), we cannot take into account any conservation law. For this reason, we do consider only global in time solutions for small data and, rather than working with Strichartz estimates, we carry on our analysis in the classical energy space by means of the aforementioned group Fourier transform.

Throughout the paper we denote by $L^q(G)$ the space of $q$-summable functions with respect to the normalized Haar measure on $G$ for $1 \leq q < \infty$ (respectively, essentially bounded for $q = \infty$) and for $s > 0$ and $q \in (1, \infty)$ the Sobolev space $H^{s,q}(G)$ is defined as the space

$$H^{s,q}(G) := \left\{ f \in L^q(G) : (-\mathcal{L})^{s/2} f \in L^q(G) \right\}$$

equipped with the norm $\| f \|_{H^{s,q}(G)} := \| f \|_{L^q(G)} + \| (-\mathcal{L})^{s/2} f \|_{L^q(G)}$. As customary, the Hilbert space $H^{s,2}(G)$ is simply denoted by $H^s(G)$.

Let us state the main result of this work, which is the global in time existence of small data solutions for the semilinear Cauchy problem (1).

**Theorem 1** Let $G$ be a compact, connected Lie group with topological dimension $n$ such that $n \geq 3$. Let $p > 1$ such that $p \leq \frac{n}{n-2}$. Then, there exists $\varepsilon_0 > 0$ such that for any $(u_0, u_1) \in H^1(\mathbb{R}) \times L^2(\mathbb{R})$ such that

$$\|(u_0, u_1)\|_{H^1(\mathbb{R}) \times L^2(\mathbb{R})} \leq \varepsilon_0$$

the semilinear Cauchy problem (1) admits a uniquely determined mild solution

$$u \in C\left([0, \infty), H^1(\mathbb{R})\right) \cap C^1\left([0, \infty), L^2(\mathbb{R})\right).$$

Furthermore, $u$ satisfies the decay estimates

$$\|u(t, \cdot)\|_{L^2(\mathbb{R})} \leq C d_{b,m^2}(t) \|(u_0, u_1)\|_{H^1(\mathbb{R}) \times L^2(\mathbb{R})},$$

$$\|(-\mathcal{L})^{1/2} u(t, \cdot)\|_{L^2(\mathbb{R})} \leq C d_{b,m^2}(t) \|(u_0, u_1)\|_{H^1(\mathbb{R}) \times L^2(\mathbb{R})},$$

$$\|\partial_t u(t, \cdot)\|_{L^2(\mathbb{R})} \leq C d_{b,m^2}(t) \|(u_0, u_1)\|_{H^1(\mathbb{R}) \times L^2(\mathbb{R})},$$

for any $t \geq 0$, where $C$ is a positive constant and the decay function $d_{b,m^2}(t)$ is given by

$$d_{b,m^2}(t) = \begin{cases} 
 e^{-\frac{b}{2}t} & \text{if } b^2 < 4m^2, \\
 t e^{-\frac{b}{2}t} & \text{if } b^2 = 4m^2, \\
 e\left(-\frac{b}{2}+\sqrt{\frac{b^2}{4}-m^2}\right)t & \text{if } b^2 > 4m^2.
\end{cases}$$
Remark 1 The upper bound for the exponent \( p \) in Theorem 1 is required in order to apply a Gagliardo–Nirenberg type inequality proved in [30, Remark 1.7]. The restriction \( n \geq 3 \) on the topological dimension is made to fulfill the conditions for the application of this inequality too and it can be removed by looking for solutions in spaces with less regularity, namely, in \( C([0, \infty), H^s_L(\mathbb{G})) \cap C^1([0, \infty), L^2(\mathbb{G})) \) for some \( s \in (0, 1) \).

Remark 2 Let us stress that the decay estimates for the solution of the semilinear Cauchy problem are exactly the same ones as for the solution of the corresponding linear homogeneous problem (cf. Proposition 1). This is a consequence of the construction of the family of evolution spaces \( \{X(T)\}_{T>0} \) equipped with a suitable weighted norm in the proof of Theorem 1 (cf. (6) in Sect. 2).

Remark 3 Note that in the statement of Theorem 1 no further lower bound for \( p > 1 \) is required. This remarkable property is due to the exponential decay rates in the \( L^2(\mathbb{G}) \)–\( L^2(\mathbb{G}) \) estimates for the solution to the corresponding linear homogeneous Cauchy problem from Proposition 1 below.

Notations

Throughout this paper the following notations are used: as in the introduction, \( \mathcal{L} \) denotes the Laplace–Beltrami operator on \( \mathbb{G} \); \( \text{Tr}(A) = \sum_{j=1}^d a_{jj} \) denotes the trace of the matrix \( A = (a_{ij})_{1 \leq i,j \leq d} \in \mathbb{C}^{d \times d} \) while \( A^* = (\overline{a_{ij}})_{1 \leq i,j \leq d} \) denotes its adjoint matrix; \( I_d \in \mathbb{C}^{d \times d} \) is the identity matrix; \( dx \) stands for the normalized Haar measure on the compact group \( \mathbb{G} \); finally, we write \( f \lesssim g \) when there exists a positive constant \( C \) such that \( f \leq Cg \).

2 Philosophy of Our Approach

In this section we will clarify our strategy to prove Theorem 1. First, we provide the notion of mild solutions to (1) in our framework. By applying Duhamel’s principle, the solution to the linear problem

\[
\begin{align*}
\begin{cases}
\partial_t^2 u - \mathcal{L} u + b \partial_t u + m^2 u &= F(t, x), & x \in \mathbb{G}, \ t > 0, \\
u(0, x) &= u_0(x), & x \in \mathbb{G}, \\
\partial_t u(0, x) &= u_1(x), & x \in \mathbb{G}
\end{cases}
\end{align*}
\]

(5)

can be represented as

\[
u(t, x) = u_0(x) \ast(x) E_0(t, x; b, m^2) + u_1(x) \ast(x) E_1(t, x; b, m^2) + \int_0^t F(s, x) \ast(x) E_1(t - s, x; b, m^2) \, ds,
\]

where \( E_0(t, x; b, m^2) \) and \( E_1(t, x; b, m^2) \) denote, respectively, the fundamental solutions to the linear Cauchy problem (5) in the homogeneous case \( F = 0 \) with
initial data \((u_0, u_1) = (\delta_0, 0)\) and \((u_0, u_1) = (0, \delta_0)\). Let us point out that in order to get the previous representation formula we applied the invariance by time translations for the differential operator \(\partial_t^2 - L + b \partial_t + m^2\) and the property \(L(v \ast(x) E_1(t, \cdot; b, m^2)) = v \ast(x) L(E_1(t, \cdot; b, m^2))\) for any left–invariant differential operator \(L\) on \(G\).

The function \(u\) is said a mild solution to (1) on \([0, T]\) if \(u\) is a fixed point for the nonlinear integral operator \(N\) defined by

\[
N : u \in X(T) \rightarrow Nu(t, x) \doteq u_0(x) \ast(x) E_0(t, x; b, m^2) + u_1(x) \ast(x) E_1(t, x; b, m^2) + \int_0^t |u(s, x)|^p \ast(x) E_1(t - s, x; b, m^2) \, ds
\]

on the evolution space \(X(T) \doteq \mathcal{C}([0, T], H^1_{\mathcal{L}}(G)) \cap \mathcal{C}([0, T], L^2(G))\), endowed with the norm

\[
\|u\|_{X(T)} \doteq \sup_{t \in [0, T]} \left( \|d_{b, m^2}(t)\|^{-1} \left( \|u(t, \cdot)\|_{L^2(G)} + \|(-\mathcal{L})^{1/2} u(t, \cdot)\|_{L^2(G)} + \|\partial_t u(t, \cdot)\|_{L^2(G)} \right) \right).
\] (6)

If we show the validity of the inequalities

\[
\|Nu\|_{X(T)} \leq C\|(u_0, u_1)\|_{H^1_{\mathcal{L}}(G) \times L^2(G)} + C\|u\|^p_{X(T)},
\] (7)

\[
\|Nu - Nv\|_{X(T)} \leq C\|u - v\|_{X(T)} \left( \|u\|_{X(T)}^{p-1} + \|v\|_{X(T)}^{p-1} \right),
\] (8)

for any \(u, v \in X(T)\) and for a suitable constant \(C > 0\) independent of \(T\), then, by Banach’s fixed point theorem it follows that \(N\) admits a uniquely determined fixed point \(u\) provided that \(\|(u_0, u_1)\|_{H^1_{\mathcal{L}}(G) \times L^2(G)}\) is small enough. This function \(u\) is our mild solution to (1) on \([0, T]\). Furthermore, thanks to the fact that (7) and (8) hold uniformly with respect to \(T\), this solution can be prolonged and defined for any \(t \in (0, \infty)\).

Before considering the semilinear Cauchy problem, we will determine \(L^2(G) – L^2(G)\) estimates for the solution to (2) via the group Fourier transform with respect to the spatial variable. We point out explicitly that the definition of the norm on the space \(X(T)\) in (6) is compatible with the decay estimates for the solution to the linear Cauchy problem (2) that we will determine afterwards (see also [8,Section 18.1.1] for a detailed overview on this topic). After that these estimates will have been established, we could show the global existence of small data solutions for (1) by applying a Gagliardo–Nirenberg type inequality derived recently in [30] (cf. Lemma 1 in Sect. 5).

The remaining part of the paper is organized as follows: in Sect. 3 we recall the main tools from Fourier Analysis and representation theory on compact Lie groups which are necessary for our approach; in Sect. 4, we establish \(L^2(G) – L^2(G)\) estimates for the solution of (2) and its first order derivatives; in Sect. 5 it will be shown that
the operator \( N \) is a contraction on \( X(T) \) provided that the data are sufficiently small in the energy space; finally, we provide some concluding remarks in Sect. 6.

3 Group Fourier Transform

Let us recall some results on Fourier Analysis on compact Lie groups analogously as in [25, Sect. 2.1]. For further details on this topic we refer to the monographs [9, 29].

A continuous unitary representation \( \xi : G \to \mathbb{C}^{d_\xi} \times \mathbb{C}^{d_\xi} \) of dimension \( d_\xi \) is a continuous group homomorphism from \( G \) to the group of unitary matrix \( U(d_\xi, \mathbb{C}) \), that is, \( \xi(xy) = \xi(x)\xi(y) \) and \( \xi(x)^* = \xi(x)^{-1} \) for all \( x, y \in G \) and the elements \( \xi_{ij} : G \to \mathbb{C} \) of the matrix representation \( \xi \) are continuous functions for all \( i, j \in \{1, \ldots, d_\xi\} \). Two representations \( \xi, \eta \) of \( G \) are said equivalent if there exists an invertible intertwining operator \( A \) such that \( A\xi(x) = \eta(x)A \) for any \( x \in G \). A subspace \( W \subset \mathbb{C}^{d_\xi} \) is \( \xi \)-invariant if \( \xi(x) \cdot W \subset W \) for any \( x \in G \). A representation \( \xi \) is irreducible if the only \( \xi \)-invariant subspaces are the trivial ones \( \{0\}, \mathbb{C}^{d_\xi} \).

The unitary dual \( \hat{G} \) of the compact Lie group \( G \) consists of the equivalence class \([\xi]\) of continuous irreducible unitary representation \( \xi : G \to \mathbb{C}^{d_\xi} \times \mathbb{C}^{d_\xi} \).

Given \( f \in L^1(G) \), its Fourier coefficients at \([\xi]\) in \( \hat{G} \) is defined by

\[
\hat{f}(\xi) = \int_G f(x)\xi(x)^*dx \in \mathbb{C}^{d_\xi} \times \mathbb{C}^{d_\xi},
\]

where the integral is taken with respect to the Haar measure on \( G \).

If \( f \in L^2(G) \), then, the Fourier series representation for \( f \) is given by

\[
f(x) = \sum_{[\xi]\in\hat{G}} d_\xi \text{Tr} (\xi(x)\hat{f}),
\]

where hereafter just one irreducible unitary matrix representation is picked in the sum for each equivalence class \([\xi]\) in \( \hat{G} \). Furthermore, for \( f \in L^2(G) \) Plancherel formula takes the following form

\[
\|f\|_{L^2(G)}^2 = \sum_{[\xi]\in\hat{G}} d_\xi \|\hat{f}(\xi)\|_{HS}^2,
\]

where the Hilbert – Schmidt norm of the matrix \( \hat{f}(\xi) \) is defined as follows:

\[
\|\hat{f}(\xi)\|_{HS}^2 = \text{Tr} \left( \hat{f}(\xi)\hat{f}(\xi)^* \right) = \sum_{i,j=1}^{d_\xi} |\hat{f}(\xi)_{ij}|^2.
\]

For our analysis it is important to understand the behavior of the group Fourier transform with respect to the Laplace–Beltrami operator \( \mathcal{L} \). Given \([\xi]\) in \( \hat{G} \), then, all
\( \xi_{ij} \) are eigenfunctions for \( \mathcal{L} \) with the same not positive eigenvalue \( -\lambda_{\xi}^2 \), namely,

\[
-\mathcal{L}\xi_{ij}(x) = \lambda_{\xi}^2 \xi_{ij}(x) \quad \text{for any } x \in \mathbb{G} \text{ and for all } i, j \in \{1, \ldots, d_{\xi}\}.
\]

This means that the symbol of \( \mathcal{L} \) is

\[
\sigma_{\mathcal{L}}(\xi) = -\lambda_{\xi}^2 I_d_{\xi}, \tag{10}
\]

that is, \( \widehat{\mathcal{L}f}(\xi) = \sigma_{\mathcal{L}}(\xi) \hat{f}(\xi) = -\lambda_{\xi}^2 \hat{f}(\xi) \) for any \( \xi \in \mathbb{G} \).

Finally, through Plancherel formula for \( s > 0 \) we have

\[
\|f\|^2_{H^s_{\mathcal{L}}(\mathbb{G})} = \|(-\mathcal{L})^{s/2}f\|_{L^2(\mathbb{G})}^2 = \sum_{[\xi] \in \mathbb{G}} d_{\xi} \lambda_{\xi}^{2s} \|\hat{f}(\xi)\|^2_{HS}.
\]

## 4 \( L^2(\mathbb{G})-L^2(\mathbb{G}) \) Estimates for the Solution to the Linear Homogeneous Problem

In this section, we derive \( L^2(\mathbb{G})-L^2(\mathbb{G}) \) estimates for the solution to the linear Cauchy problem (2). We follow the approach from [10], which have been recently applied to study other semilinear hyperbolic models in compact Lie groups in [25, 26]. The main idea of this approach is to employ the group Fourier transform with respect to the spatial variable \( x \) to get an explicit expression for the \( L^2(\mathbb{G}) \) norms of \( u(t, \cdot) \), \( (-\mathcal{L})^{1/2}u(t, \cdot) \) and \( \partial_t u(t, \cdot) \), respectively. Plancherel formula in the framework of compact Lie groups plays a fundamental role in this step.

**Proposition 1** Let \( \mathbb{G} \) be a compact Lie group. Let us assume \( u_0 \in H^1_{\mathcal{L}}(\mathbb{G}), u_1 \in L^2(\mathbb{G}) \) and let \( u \in \mathcal{C}([0, \infty), H^1_{\mathcal{L}}(\mathbb{G})) \cap \mathcal{C}^1([0, \infty), L^2(\mathbb{G})) \) be the solution to the homogeneous Cauchy problem (2).

Then, the following \( L^2(\mathbb{G})-L^2(\mathbb{G}) \) decay estimates are satisfied

\[
\|u(t, \cdot)\|_{L^2(\mathbb{G})} \leq C d_{b,m^2}(t) \left( \|u_0\|_{L^2(\mathbb{G})} + \|u_1\|_{L^2(\mathbb{G})} \right), \tag{11}
\]

\[
\|(-\mathcal{L})^{1/2}u(t, \cdot)\|_{L^2(\mathbb{G})} \leq C d_{b,m^2}(t) \left( \|u_0\|_{H^1_{\mathcal{L}}(\mathbb{G})} + \|u_1\|_{L^2(\mathbb{G})} \right), \tag{12}
\]

\[
\|\partial_t u(t, \cdot)\|_{L^2(\mathbb{G})} \leq C d_{b,m^2}(t) \left( \|u_0\|_{H^1_{\mathcal{L}}(\mathbb{G})} + \|u_1\|_{L^2(\mathbb{G})} \right), \tag{13}
\]

for any \( t \geq 0 \), where \( C \) is a positive multiplicative constant and the decay function \( d_{b,m^2} \) is defined in (4).

**Remark 4** In the threshold case \( b^2 = 4m^2 \) we may refine (12), considering \( e^{-\frac{b^2}{4}} \) as decay rate (cf. (23) in the proof below). However, in the proof of the global existence result the decay rate in (12) suffices to apply Banach’s fixed point theorem for any \( p > 1 \) without requiring a further lower bound for \( p \).

**Remark 5** We point out that no connectedness is required for \( \mathbb{G} \) in the study of the linear Cauchy problem (2) in the statement of Proposition 1.
Representation Formula for the Fourier Coefficient \( \hat{u}(t, \xi)_{k\ell} \)

Let \( u \) solve (2). By \( \hat{u}(t, \xi) = (\hat{u}(t, \xi)_{k\ell})_{1 \leq k, \ell \leq d_\xi} \in \mathbb{C}^{d_\xi \times d_\xi} \), \( [\xi] \in \hat{\mathbb{G}} \) we denote the group Fourier transform of \( u \) with respect to the \( x \)-variable. Therefore, \( \hat{u}(t, \xi) \) is a solution of the Cauchy problem for the system of ODEs (with the size of the system that depends on the representation \( \xi \))

\[
\begin{align*}
\partial_t^2 \hat{u}(t, \xi) - \sigma \mathcal{L}(\xi) \hat{u}(t, \xi) + b \partial_t \hat{u}(t, \xi) + m^2 \hat{u}(t, \xi) &= 0, \quad t > 0, \\
\hat{u}(0, \xi) &= \hat{u}_0(\xi), \\
\partial_t \hat{u}(0, \xi) &= \hat{u}_1(\xi).
\end{align*}
\]

By using the symbol of the Laplace–Beltrami operator in (10), we obtain that the previous system is actually decoupled in \( d_\xi^2 \) independent scalar ODEs, namely,

\[
\begin{align*}
\partial_t^2 \hat{u}(t, \xi)_{k\ell} + \lambda^2_{k\ell} \hat{u}(t, \xi)_{k\ell} + b \partial_t \hat{u}(t, \xi)_{k\ell} + m^2 \hat{u}(t, \xi)_{k\ell} &= 0, \quad t > 0, \\
\hat{u}(0, \xi)_{k\ell} &= \hat{u}_0(\xi)_{k\ell}, \\
\partial_t \hat{u}(0, \xi)_{k\ell} &= \hat{u}_1(\xi)_{k\ell},
\end{align*}
\]

for any \( k, \ell \in \{1, \ldots, d_\xi\} \). Straightforward computations lead to the representation formula

\[
\hat{u}(t, \xi)_{k\ell} = e^{-\frac{b}{2} t} G_0(t; b, m^2; \xi) \hat{u}_0(\xi)_{k\ell} + e^{-\frac{b}{2} t} G_1(t; b, m^2; \xi) \hat{u}_1(\xi)_{k\ell}
\]

for the solution to the linear homogeneous Cauchy problem (14), where

\[
G_0(t; b, m^2; \xi) \doteq \begin{cases} 
\cosh \left( \frac{\sqrt{b^2/4 - m^2} - \lambda_{k\ell}}{t} \right) & \text{if } \lambda_{k\ell}^2 < \frac{b^2}{4} - m^2, \\
1 & \text{if } \lambda_{k\ell}^2 = \frac{b^2}{4} - m^2, \\
\cos \left( \frac{\sqrt{\lambda_{k\ell}^2 + m^2} - \sqrt{b^2/4}}{t} \right) & \text{if } \lambda_{k\ell}^2 > \frac{b^2}{4} - m^2,
\end{cases}
\]

\[
G_1(t; b, m^2; \xi) \doteq \begin{cases} 
\sinh \left( \frac{\sqrt{b^2/4 - m^2} - \lambda_{k\ell}}{t} \right) & \text{if } \lambda_{k\ell}^2 < \frac{b^2}{4} - m^2, \\
\sqrt{\frac{b^2}{4} - m^2} - \lambda_{k\ell} & \text{if } \lambda_{k\ell}^2 = \frac{b^2}{4} - m^2, \\
\sin \left( \frac{\sqrt{\lambda_{k\ell}^2 + m^2} - \sqrt{b^2/4}}{t} \right) & \text{if } \lambda_{k\ell}^2 > \frac{b^2}{4} - m^2.
\end{cases}
\]

Notice that \( G_0(t; b, m^2; \xi) = \partial_t G_1(t; b, m^2; \xi) \) for any \([\xi] \in \hat{\mathbb{G}}\).
Finally, by (15) for any [ξ] ∈ ̂G and any k, ℓ ∈ {1, . . . , dξ} we derive the following representation for the time derivative of ̂u(t, ξ)kℓ

\[ \partial_t ̂u(t, ξ)kℓ = e^{-b2t} G_0(t; b, m^2; ξ) ̂u_1(ξ)kℓ - e^{-b2t} G_1(t; b, m^2; ξ) \left( \frac{b2}{2} ̂u_1(ξ)kℓ + (λξ^2 + m^2) ̂u_0(ξ)kℓ \right). \] (17)

Let us prove now Proposition 1. We will consider separately the three subcases b2 \(\gg\) 4m2.

**Case b2 < 4m2**

In this case the characteristic roots are always complex conjugate and in the representation formula (15) it has sense to consider only the case \(λξ^2 > b2 - m^2\), since all eigenvalues \(\{λξ^2\}_{[ξ] ∈ ̂G}\) of \(-\mathcal{L}\) are nonnegative. So, we can estimate

\[ | ̂u(t, ξ)kℓ| \lesssim e^{-b2t} (| ̂u_0(ξ)kℓ| + | ̂u_1(ξ)kℓ|), \]

\[ λξ | ̂u(t, ξ)kℓ| \lesssim e^{-b2t} ((1 + λξ)| ̂u_0(ξ)kℓ| + | ̂u_1(ξ)kℓ|), \] (18)

for any t \(\geq 0\). Similarly, by (17) we get

\[ | ̂u(t, ξ)kℓ| \lesssim e^{-b2t} (1 + λξ)| ̂u_0(ξ)kℓ| + | ̂u_1(ξ)kℓ| \] (19)

for any t \(\geq 0\). Combining (18) and (19) and using Plancherel formula, we have

\[ \|(-\mathcal{L})^{i/2} \partial^j_t u(t, \cdot)\|^2_{L^2(\mathbb{G})} = \sum_{[ξ] ∈ ̂G} dξ \sum_{k, ℓ=1}^{dξ} λξ^{2j} | ̂u(t, ξ)kℓ|^2 \]

\[ \lesssim e^{-bt} \sum_{[ξ] ∈ ̂G} dξ \sum_{k, ℓ=1}^{dξ} \left( (1 + λξ^2)^{(i+j)} | ̂u_0(ξ)kℓ|^2 + | ̂u_1(ξ)kℓ|^2 \right) \]

\[ = e^{-bt} \left( \|u_0\|^2_{H_{\mathcal{L}}^{i+j}(\mathbb{G})} + \|u_1\|^2_{L^2(\mathbb{G})} \right) \] (20)

for any i, j \(\in\{0, 1\}\) such that 0 \(\leq i + j \leq 1\) and any t \(\geq 0\). Clearly (20) implies (11), (12) and (13) in the case b2 < 4m2 (here we used \(H_{\mathcal{L}}^0(\mathbb{G}) = L^2(\mathbb{G})\)).

**Case b2 = 4m2**

In this case the representation formula (15) is just

\[ ̂u(t, ξ)kℓ = e^{-b2t} \cos(λξt) ̂u_0(ξ)kℓ + e^{-b2t} \frac{\sin(λξt)}{λξ} \left( ̂u_1(ξ)kℓ + \frac{b}{2} ̂u_0(ξ)kℓ \right) \]
for \([\xi] \in \hat{G}\) such that \(\lambda_\xi^2 > 0\), and

\[
\hat{u}(t, \xi)_{k\ell} = e^{-\frac{b}{2}t} \hat{u}_0(\xi)_{k\ell} + t e^{-\frac{b}{2}t} \left( \hat{u}_1(\xi)_{k\ell} + \frac{b}{2} \hat{u}_0(\xi)_{k\ell} \right)
\]

for \([\xi] \in \hat{G}\) such that \(\lambda_\xi^2 = 0\).

Note that we have to consider necessarily the second case, since for the trivial 1-dimensional representation \(1 : x \in G \rightarrow 1 \in \mathbb{C}\) we have \(-L(1) = 0\). Therefore,

\[
|\hat{u}(t, \xi)_{k\ell}| \lesssim t e^{-\frac{b}{2}t} \left( |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}| \right),
\]

\[
\lambda_\xi |\hat{u}(t, \xi)_{k\ell}| \lesssim e^{-\frac{b}{2}t} \left( (1 + \lambda_\xi) |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}| \right),
\]

(21)

for any \(t \geq 0\). In the second inequality, we can drop the factor \(t\) as the spectrum of \(-L\) is discrete and has no finite cluster points.

Similarly, from (17) it follows

\[
|\partial_t \hat{u}(t, \xi)_{k\ell}| \lesssim e^{-\frac{b}{2}t} \left( (1 + \lambda_\xi) |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}| \right)
\]

(22)

for any \(t \geq 0\). Applying Plancherel formula twice as in (20) and using (21) and (22) to control the Fourier coefficients, we obtain (11), (12) and (13) in the case \(b^2 = 4m^2\).

**Case \(b^2 > 4m^2\)**

In this last case, the characteristic roots may be either complex conjugate or coincident or real distinct, depending on the range for \(\lambda_\xi^2\). Comparing all possible cases in (16), we get

\[
|\hat{u}(t, \xi)_{k\ell}| \lesssim e^{-\frac{b}{2}t + \sqrt{\frac{b^2}{4} - m^2}} t \left( |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}| \right),
\]

\[
\lambda_\xi |\hat{u}(t, \xi)_{k\ell}| \lesssim e^{-\frac{b}{2}t + \sqrt{\frac{b^2}{4} - m^2}} t \left( (1 + \lambda_\xi) |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}| \right),
\]

(23)

for any \(t \geq 0\). We point out that the regularity is provided from the case with complex conjugate characteristic roots whereas the decay rates is given by the continuous irreducible unitary representations with \(\lambda_\xi^2 = 0\).

Analogously, from (17) we get

\[
|\partial_t \hat{u}(t, \xi)_{k\ell}| \lesssim e^{-\frac{b}{2}t + \sqrt{\frac{b^2}{4} - m^2}} t \left( (1 + \lambda_\xi) |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}| \right)
\]

(24)

for any \(t \geq 0\).
Consequently, by Plancherel formula combined with (23) and (24), we find
\[
\|(-\mathcal{L}^{j/2})\partial_t^j u(t, \cdot)\|_{L^2(G)}^2 = \sum_{[\xi] \in \hat{G}} d_\xi \sum_{k, \ell = 1}^{d_\xi} \lambda^2_\xi \left| \partial_t^j \hat{u}(t, \xi)_{k\ell} \right|^2
\]
\[
\leq e^{-b + \sqrt{b^2 - 4m^2}} t \sum_{[\xi] \in \hat{G}} d_\xi \sum_{k, \ell = 1}^{d_\xi} \left( 1 + \lambda^2_\xi \right)^{(i+j)} \left| \hat{u}_0(\xi)_{k\ell} \right|^2 + \left| \hat{u}_1(\xi)_{k\ell} \right|^2
\]
\[
= e^{-b + \sqrt{b^2 - 4m^2}} t \left( \|u_0\|^2_{H^{i+j}_{\mathcal{L}}(G)} + \|u_1\|^2_{L^2(G)} \right) \tag{25}
\]
for any \(i, j \in \{0, 1\}\) such that \(0 \leq i + j \leq 1\) and any \(t \geq 0\). It is easy to see that (25) implies (11), (12) and (13) in the case \(b^2 > 4m^2\).

5 Proof of Theorem 1

A fundamental tool to prove the Theorem 1 is the following Gagliardo–Nirenberg type inequality, whose proof can be found in [30] in a more general setting (see also [25, Corollary 2.3]).

Lemma 1 Let \(G\) be a connected unimodular Lie group with topological dimension \(n \geq 3\). For any \(q \geq 2\) such that \(q \leq \frac{2n}{n-2}\) the following Gagliardo–Nirenberg type inequality holds
\[
\|f\|_{L^q(G)} \lesssim \|f\|_{H^1_{\mathcal{L}}(G)}^{\theta(n,q)} \|f\|_{L^2(G)}^{1-\theta(n,q)} \tag{26}
\]
for any \(f \in H^1_{\mathcal{L}}(G)\), where \(\theta(n,q) \doteq n \left( \frac{1}{2} - \frac{1}{q} \right)\).

As we explained in Sect. 2, in order to prove Theorem 1 it suffices to show the validity of (7) and (8).

It is helpful to rewrite \(Nu = u^\ln + Ju\), where
\[
u^\ln(t, x) \doteq u_0(x) \ast_{(x)} E_0(t, x; b, m^2) + u_1(x) \ast_{(x)} E_1(t, x; b, m^2),
\]
\[Ju(t, x) \doteq \int_0^t |u(s, x)|^p \ast_{(x)} E_1(t - s, x; b, m^2) \, ds. \tag{27}\]

Let us get started by estimating \(\|Nu\|_{X(T)}\) for \(u \in X(T)\). By Proposition 1 it results
\[
\|u^\ln\|_{X(T)} \lesssim \|(u_0, u_1)\|_{H^1_{\mathcal{L}}(G) \times L^2(G)}. \tag{28}\]

We recall from (27) the definition of the term \(Ju\). In the next step, we are going to use the invariance by time translations of (2) that allows us to estimate immediately the
term $|u(s, x)|^p \ast_{(x)} E_1(t - s, x; b, m^2)$ by using Proposition 1. Hence, combining the Gagliardo–Nirenberg type inequality from Lemma 1 and (6), we get

$$\|d_t^i (-\mathcal{L})^{j/2} J u(t, \cdot)\|_{L^2(\mathbb{G})} \lesssim \int_0^t d_{b, m^2}(t - s) \|u(s, \cdot)\|^p_{L^2_p(\mathbb{G})} \, ds$$

$$\lesssim \int_0^t d_{b, m^2}(t - s) \|u(s, \cdot)\|^{p\theta(n,2p)}_{H^1(\mathbb{G})} \|u(s, \cdot)\|^{p(1-\theta(n,2p))}_{L^2(\mathbb{G})} \, ds$$

$$\lesssim \int_0^t d_{b, m^2}(t - s)(d_{b, m^2}(s))^p \, ds \|u\|^p_{X(t)}$$

$$\lesssim d_{b, m^2}(t) \|u\|^p_{X(t)}$$

(29)

for $i, j \in \{0, 1\}$ such that $0 \leq i + j \leq 1$. We underline that the employment of (26) in the second step of the previous chain of inequality, we have to require the condition $p \leq \frac{n}{n-2}$ in Theorem 1. So, from (28) and (29) we obtain (7).

Let us derive now (8). We point out that $J u - J v$ is a solution of the equation with homogeneous Cauchy data and with source term $|u|^p - |v|^p$. Combining Hölder’s inequality and the inequality

$$||u|^p - |v|^p| \leq p|u - v|(|u|^{p-1} + |v|^{p-1}),$$

which is a straightforward consequence of the mean value theorem applied to the $p$-power function, we get

$$\|\|u(s, \cdot)|^p - |v(s, \cdot)|^p\|_{L^2(\mathbb{G})}$$

$$\lesssim \|u(s, \cdot) - v(s, \cdot)\|_{L^2_p(\mathbb{G})} \left(\|u(s, \cdot)\|^{p-1}_{L^2_p(\mathbb{G})} + \|v(s, \cdot)\|^{p-1}_{L^2_p(\mathbb{G})}\right).$$

Applying the previous inequality and (26) to estimate each $L^2_p(\mathbb{G})$-norm that appears on the right-hand side, we find

$$\|d_t^i (-\mathcal{L})^{j/2}(J u(t, \cdot) - J v(t, \cdot))\|_{L^2(\mathbb{G})}$$

$$\lesssim \int_0^t d_{b, m^2}(t - s) \|u(s, \cdot)\|^p_{L^2(\mathbb{G})} \, ds$$

$$\lesssim \int_0^t d_{b, m^2}(t - s)(d_{b, m^2}(s))^p \, ds \|u - v\|_{X(t)} \left(\|u\|^{p-1}_{X(t)} + \|v\|^{p-1}_{X(t)}\right)$$

$$\lesssim d_{b, m^2}(t) \|u - v\|_{X(t)} \left(\|u\|^{p-1}_{X(t)} + \|v\|^{p-1}_{X(t)}\right)$$

(30)

for $i, j \in \{0, 1\}$ such that $0 \leq i + j \leq 1$. From (6) and (30) it follows (8).

Note that thanks to the exponential decay rate $d_{b, m^2}(t)$ both in (29) and (30) we have the uniform boundedness of the integral.
\[
(d_{b,m^2}(t))^{-1} \int_0^t \left( d_{b,m^2}(t-s)(d_{b,m^2}(s))^p \right) ds
\]  
(31)

with respect to \( t \) without requiring further conditions on \( p \).

### 6 Final Remarks

In [25, 26] it is emphasized how the global dimension of \( G \) (which is 0 in the compact case) has a relevant role in the search for globally defined solution for the semilinear Cauchy problem with power nonlinearity associated to the damped wave operator \( \partial_t^2 - L + \partial_t \) and to the wave operator \( \partial_t^2 - L \), respectively. In both cases, we may not prove any global existence result. On the contrary, for any \( p > 1 \) local in time solutions to these semilinear problem blow up in finite time under suitable sign assumption for the Cauchy data. In this paper, we showed that the contemporary presence of a damping term and of a mass term reverses completely the situation.

As we mentioned in the introduction, this fact has been already observed in the Euclidean case (also in the case with fractional in space operators [7] or with suitable time-dependent coefficients for the damping and mass terms [15]) and for graded Lie group. Nevertheless, it is remarkable to observe this phenomenon even in compact (and connected) Lie groups in consideration of what remarked above for the damped wave operator and for the wave operator, for which the corresponding semilinear Cauchy problems admit a finite critical exponent in the Euclidean case (Fujita exponent [18, 23, 34, 37] and Strauss exponent [12, 16, 17, 19, 22, 31, 32, 36, 38], respectively) or in the special case of the damped wave equation in the Heisenberg group (see [11, 24]).

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