Understanding the Limitations of Gyrochronology for Old Field Stars

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Abstract

Nearly half a century has passed since the initial indications that stellar rotation slows while chromospheric activity weakens with a power-law dependence on age, the so-called Skumanich relations. Subsequent characterization of the mass-dependence of this behavior up to the age of the Sun led to the advent of gyrochronology, which uses the rotation rate of a star to infer its age from an empirical calibration. The efficacy of the method relies on predictable angular momentum loss from a stellar wind entrained in the large-scale magnetic field produced by global dynamo action. Recent observational evidence suggests that the global dynamo begins to shut down near the middle of a star’s main-sequence lifetime, leading to a disruption in the production of large-scale magnetic field, a dramatic reduction in angular momentum loss, and a breakdown of gyrochronology relations. For solar-type stars this transition appears to occur near the age of the Sun, when rotation becomes too slow to imprint Coriolis forces on the global convective patterns, reducing the shear induced by differential rotation, and disrupting the large-scale dynamo. We use data from Barnes to reveal the signature of this transition in the observations that were originally used to validate gyrochronology. We propose that chromospheric activity may ultimately provide a more reliable age indicator for older stars, and we suggest that asteroseismology can be used to help calibrate activity–age relations for field stars beyond the middle of their main-sequence lifetimes.

Key words: stars: activity – stars: evolution – stars: magnetic field – stars: rotation – stars: solar-type

1. Background

Stars are born with a range of initial rotation rates and magnetic field strengths, and beyond the saturated regime the two properties are intricately linked for as long as a global dynamo continues to operate. The large-scale magnetic field gradually slows the rotation over time (e.g., see Réville et al. 2015; Garaffo et al. 2016). Through a process known as magnetic braking, charged particles in the stellar wind follow the magnetic field lines out to the Alfvén radius, shedding angular momentum in the process. In turn, nonuniform rotation modifies the morphology of the magnetic field (e.g., see Brown et al. 2010). Solar-like differential rotation, with a faster equator and slower poles, is a natural consequence of convection in the presence of substantial Coriolis forces (Miesch 2005). The resulting shear wraps up the large-scale poloidal field into a toroidal configuration that ultimately leads to the emergence of active regions on the surface. Through these basic physical processes, stellar rotation and magnetism diminish together over time, each feeding off the other. The mutual feedback can continue as long as rotation and magnetism are coupled through a global dynamo.4

Nearly half a century ago, Skumanich (1972) planted the observational seeds of this consensus view of magnetic stellar evolution. Both the theoretical foundations and the constraints from young clusters improved steadily over the intervening decades (e.g., see Soderblom et al. 1993, and references therein). But the Sun remained the oldest calibrator, so the empirical relations were largely untested beyond stellar middle-age. Barnes (2007) put forward a more quantitative formulation of the rotation-age relation (so-called gyrochronology), establishing the mass-dependence of stellar spin-down from observations of young clusters and using the Sun to determine the age-dependence. Given only the $B-V$ color and rotation period ($P_{\text{rot}}$) of a star, gyrochronology yielded an empirical stellar age with a precision of $15\%–20\%$. Barnes (2010) revised this formulation to account for varying initial conditions ($P_{\text{ip}}$, important in young clusters), and to map the mass-dependence onto a convective turnover time ($\tau_c$) derived from the stellar models of Barnes & Kim (2010). This approach more faithfully reproduced the distribution of rotation periods in young clusters, while yielding ages compatible with Barnes (2007) for more evolved stars.

Observations from the Kepler mission provided the first tests of gyrochronology for older clusters and for field stars beyond the age of the Sun. Meibom et al. (2011) found good agreement with expectations for the 1 Gyr cluster NGC 6811, and Meibom et al. (2015) extended this success to 2.5 Gyr with observations of the cluster NGC 6819. The first indications of unexpected behavior were uncovered by Angus et al. (2015), who found that no single gyrochronology relation could simultaneously explain the cluster data and the asteroseismic ages for old Kepler field stars with measured rotation periods. van Saders et al. (2016) confirmed anomalously fast rotation among the best characterized Kepler asteroseismic targets, and proposed a model that could explain the observations with significantly weakened magnetic braking beyond the middle of a star’s main-sequence lifetime. Metcalfe et al. (2016) found the magnetic counterpart of this rotational transition in chromospheric activity measurements of the Kepler targets, showing empirically that the activity level continues to decrease while the rotation rate remains almost constant. They suggested that the transition might be triggered by a change in the character of differential rotation that was expected from global convection simulations (Gastine et al. 2014; Brun et al. 2017).

4 By “global dynamo” we mean the mechanism that generates large-scale magnetic field, as opposed to a “local dynamo,” which may generate field on smaller scales.
Metcalfe & van Saders (2017) identified a coincident shift in stellar cycle properties, with the cycle period growing longer and the amplitude becoming weaker at nearly constant rotation.

These developments suggest a revised picture of the late stages of magnetic stellar evolution, in which the disruption of differential rotation in the absence of substantial Coriolis forces leads to a gradual decrease in the production of large-scale differential rotation in the absence of substantial Coriolis forces long before the subgiant phase. This scenario would also explain the long-period edge found by McQuillan et al. in the distribution of rotation periods with age. In Section 4, we discuss asteroseismic analogs of the Mount Wilson stars that show the largest decoupling of rotation and magnetism.

Figure 1. Difference between the chromospheric age and the gyro age, in units of the uncertainties of the chromospheric (\sigma_{\text{chromo}}) and gyro (\sigma_{\text{gyro}}) age estimates. The stars are summarized in Table 1 for numerical values of the age discrepancy for the outliers at low activity levels.

In an effort to address any skepticism about the existence of this transition, we identify its manifestation among the most evolved dwarfs in the Mount Wilson sample that were originally used to validate gyrochronology (Section 2). We then search within the Kepler asteroseismic sample to identify analogs of the Mount Wilson stars that show the largest inconsistencies between gyrochronology and chromospheric ages (Section 3), allowing us to characterize more precisely the decoupling of rotation and magnetism. In Section 4, we discuss future observations of these Mount Wilson stars with the Transiting Exoplanet Survey Satellite (TESS), and we predict that their asteroseismic ages will significantly exceed those expected from gyrochronology. We conclude in Section 5 with a discussion of the potential for reliable chromospheric ages of older stars, using asteroseismology to recalibrate the activity–age relation.

2. Gyrochronology Sample

After calibrating gyrochronology with young clusters and the Sun, Barnes (2007, hereafter B07) attempted to validate the method using a sample of bright field stars observed for decades by the Mount Wilson HK project (Wilson 1968). For the subset of 71 stars that were not known to be significantly evolved, B07 compiled B − V colors, rotation periods \(P_\text{rot}\), and mean chromospheric activity levels \(\langle R'_{\text{HK}} \rangle\) from the literature (Noyes et al. 1984; Baliunas et al. 1996; Donahue et al. 1996). B07 then used the mean activity levels to calculate chromospheric ages from the activity–age relation of Donahue (1998), which could be compared to the ages and uncertainties from gyrochronology (see Table 3 of B07). Although Donahue (1998) did not provide a method for assessing uncertainties on the calculated ages, B07 noted that discrepancies in the age estimates for members of wide binaries and triple systems suggested a mean fractional error of 46%, about 3 times the typical age uncertainties from gyrochronology. In summarizing the comparison between the two age estimates, B07 noted in the abstract: “Gyro ages for the Mount Wilson stars are shown to be in good agreement with chromospheric ages for all but the bluest stars.” Below, we use the data from Table 3 of B07 to reassess this comparison.

In Figure 1, we show the difference between the chromospheric and gyro ages, in units of the age uncertainty, plotted against the mean chromospheric activity level for 70 of the Mount Wilson stars tabulated in B07 (we omit only the M dwarf HD 95735, which has a spurious chromospheric age of 20 Gyr). On the left axis, the inconsistency between the age estimates is shown in units of the tabulated uncertainty on the gyro age, \(\sigma_{\text{gyro}}\). On the right axis, the values are scaled to reflect the larger uncertainty on the chromospheric age, \(\sigma_{\text{chromo}}\). Colored symbols separate the sample into hotter stars with \(B − V < 0.6\) (blue triangles), solar-type stars with \(0.6 < B − V < 0.8\) (yellow circles), and cooler stars with \(B − V > 0.8\) (red squares). The solar values from B07 are indicated with the symbol. As noted by B07: “apart from a slight tendency toward shorter gyro ages...there is general agreement between the chromospheric and gyro ages for this sample.” With few exceptions, the two age estimates tend to agree between activity levels of \(-4.3\) and \(-5.0\), in a band of uncertainty that stretches from \(-1\) to \(+3\) \(\sigma_{\text{chromo}}\) (\(-3\) to \(+9\sigma_{\text{gyro}}\)).

Given the relative precision of the two methods, we can assume that most of the scatter is due to the chromospheric age uncertainties. Although the median age inconsistency is indeed slightly higher for the blue stars \((B − V < 0.6)\), the most significant outliers\(^5\) are found at the lowest activity levels (\(\log R'_{\text{HK}} < -5\)). These stars are labeled with their HD numbers in Figure 1, and their properties are listed in Table 1. There is reason to be skeptical of the chromospheric ages for these Mount Wilson stars. As pointed out by B07, the chromospheric ages for these stars exceed the main-sequence.

\(^5\) Exceptions include the two Jovian exoplanet host stars GJ 504 (Kuzuhara et al. 2013) and τ Boo (Walker et al. 2008), where the stellar rotation, activity, or both could plausibly be affected by interactions with the planet.

\(^6\) Note that the identification of these stars as the largest outliers does not depend on the B07 formulation of gyrochronology. Updating all of the gyro ages to those produced by the (Barnes 2010, hereafter B10) formulation (and adopting the uncertainties from B07, because there is no prescription for calculating uncertainties in B10) yields the same conclusion.
lifetime of a typical F-type star. This argument was used by
B07 as justification for giving preference to the gyro ages.
However, even if the chromospheric ages are substantially
overestimated, the gyro ages are not necessarily correct.

Setting aside questions about the absolute reliability of
chromospheric ages, the activity levels certainly suggest that
these stars might be significantly evolved (e.g., see
Wright 2004). In the bottom half of Table 1, we provide three
additional lines of evidence that support this general conclu-
sion. First, we list spectroscopic parameters ($T_{\text{eff}}$, log $g$,
[Fe/H]) from Boeche & Grebel (2016). Barnes et al. (2016)
argue that gyrochronology has only been calibrated for dwarf
stars near solar-metallicity. They classify as subgiants any star
with log $g < 4.2$, and remove from consideration metal-poor
stars. By these definitions, all of these Mount Wilson stars
would be considered subgiants, and one of them (HD 45067)
should be discarded by virtue of its low metallicity. Second, the
luminosities shown in Table 1 (Gaia Collaboration et al. 2018)
are substantially above the main-sequence luminosity for
F-type stars (which is typically less than $2L_{\odot}$). Third, after
25 years of monitoring their chromospheric activity, Baliunas
et al. (1995) classified most of these stars as “Flat” (i.e.,
showing constant activity with fractional variations less than
1.5%) or “Long” (i.e., showing potential variability on a
timescale longer than 25 years), suggesting that their global
dynamos may have already started to shut down (Metcalfe &
van Saders 2017). The one exception is a possible 5.4 yr cycle
in HD 187691, which was assigned a false-alarm probability
grade of “Fair” by Baliunas et al. (1995). Note that the two
solar-type stars near HD 212754 in Figure 1 (HD 143761 and
HD 187691), which was assigned a false-alarm probability
grade of “Flat” by Baliunas et al. (1995). Although
there may be substantial problems with chromospheric ages at
these low activity levels (e.g., see Mamajek & Hillenbrand
2008), the corroborating evidence in Table 1 of significant
evolution suggests that gyro ages may also suffer from
systematic errors in this regime.

In the context of our revised picture of magnetic stellar
evolution (Section 1), how can we understand this inconsis-
tency between the chromospheric and gyro ages for the
Mount Wilson stars in Table 1? As discussed by van Saders
et al. (2016), the shutdown of magnetic braking appears to
occur at a critical value of the Rossby number ($\text{Ro} \equiv P_{\text{rot}}/\tau_{c}$),
the ratio of the rotation period to the convective turnover time.
Hotter stars have shallower convection zones with shorter
turnover times, so they reach the critical Rossby number at
earlier absolute ages while their rotation periods are still
relatively short. The result of this transition is a decoupling of
rotation and magnetism, with the rotation period remaining
almost constant while the chromospheric activity continues to
decline with age (Metcalfe et al. 2016). As suggested by
Metcalfe & van Saders (2017), stellar cycles appear to grow
longer and decrease their amplitude during this transition before disappearing entirely or becoming undetectable, leading
to classifications of “Long” or “Flat” in stellar cycle surveys.
With the rotation period essentially fixed, the gyro age is a
lower limit that actually reflects the age when the star stopped
spinning down in the middle of its main-sequence lifetime.

### 3. Analogs Observed by Kepler

Although asteroseismic observations do not yet exist for the
Mount Wilson stars in Table 1 (see Section 4), we can search
for analogs of these stars within the sample of Kepler
asteroseismic targets. There are currently 18 Kepler targets
with detailed asteroseismic modeling (for precise ages) that
also have known rotation periods and measured chromospheric
activity (see Metcalfe et al. 2016, their Table 1). Among these
stars, only three fall within the same range of rotation periods
and $B - V$ colors as the Mount Wilson stars in Table 1. We list
the properties of these analogs in the first three columns of
Table 2, ordered by activity level.

The chromospheric activity levels for these analogs span the
range where the magnetic transition summarized in Section 1 is
expected to occur. The critical Rossby number found by van
Saders et al. (2016) can also be understood as a critical activity
level, because the two properties are strongly correlated
(Mamajek & Hillenbrand 2008). The critical activity level is
around $-4.95$ (Brandenburg et al. 2017), so the gap between
gyrochronology and other age estimates is expected to grow
wider as stars continue to evolve to lower activity levels. We
use the $B - V$ color and $P_{\text{rot}}$ to calculate gyro ages and
uncertainties using the B07 formulation, and we use log $R_{\text{HK}}$
per chromospheric ages following Donahue (1998).
For KIC 9139151, which is just reaching the critical activity
level where rotation and magnetism are expected to decouple,
the gyro age agrees with the asteroseismic age. For the more

| $B - V$ | HD 45067 | HD 89744 | HD 107213 | HD 187691 | HD 212754 |
|---------|----------|----------|------------|------------|------------|
| 0.56    | 0.54     | 0.50     | 0.55       | 0.52       |            |
| $P_{\text{rot}}$ [day] | 8         | 9         | 9          | 10         | 12         |
| log ([$R_{\text{HK}}$]) | $-5.094$  | $-5.120$  | $-5.103$   | $-5.026$   | $-5.073$   |
| $t_{\text{gyro}}$ [Gyr] | $0.76 \pm 0.12$ | $1.11 \pm 0.19$ | $1.63 \pm 0.33$ | $1.25 \pm 0.21$ | $2.30 \pm 0.44$ |
| $t_{\text{chrono}}$ [Gyr] | 7.733 | 8.421 | 7.966 | 6.128 | 7.207 |
| $T_{\text{eff}}$ [K] | 5973 | 6149 | 6249 | 6059 | 6210 |
| log g | 3.88 | 3.94 | 4.13 | 4.06 | 3.88 |
| [Fe/H] | $-0.19$ | $+0.08$ | $+0.16$ | $+0.04$ | $-0.05$ |
| $L/L_{\odot}$ | $4.18 \pm 0.02$ | $6.38 \pm 0.02$ | $5.31 \pm 0.02$ | $2.95 \pm 0.01$ | $6.72 \pm 0.05$ |
| MWO $P_{\text{rot}}$ [yr] | Flat | Flat | Long | 5.4 (Fair) | Long |
evolved dwarfs KIC 12009504 and KIC 10963065, the gap between the gyro age and the asteroseismic age is significant. The gyro age for these stars appears to indicate the point at which a magnetic braking became inefficient and rotation stopped evolving substantially. The available data from Kepler suggests that gyrochronology is an unreliable age indicator for hotter stars beyond $\sim$2–3 Gyr (van Saders et al. 2016).

The chromospheric ages for the analogs appear to be substantially overestimated at these low activity levels, just as for the Mount Wilson stars in Table 1. The other indicators of evolutionary status that are listed in the first three columns of Table 2 generally fall within the range where gyrochronology has been calibrated. The surface gravities are all above the cut $(\log g > 4.2)$ suggested by Barnes et al. (2016), only one of the stars (KIC 10963065) is significantly metal-poor, and the Gaia luminosities are well below those of the Mount Wilson stars. Most significantly, there is no apparent reason to expect gyrochronology to fail for KIC 12009504, but both the B07 (1.06 Gyr) and B10 (1.10 Gyr) ages are wildly inconsistent with the asteroseismic age (modeling results vary between 3.10 and 4.12 Gyr, Silva Aguirre et al. 2017). If this star is actually too evolved for gyrochronology to be reliable, the rotation period should have slowed as it expanded into a subgiant, biasing the gyro age older and reducing the inconsistency with asteroseismology.

4. Predictions for TESS

The TESS mission launched successfully in April 2018, and it is expected to gather short-cadence photometry (2 minute sampling) for all of the Mount Wilson stars listed in Table 1. Detections of solar-like oscillations comparable to what was achieved by Kepler are expected in TESS targets that are $\sim$5 mag brighter (Campante et al. 2016). The minimum dwell time on each TESS sector is 27 days, comparable to the 30 day time series obtained for Kepler targets during the asteroseismic survey that was conducted in the first year (Chaplin et al. 2011b). The amplitude of solar-like oscillations scales with the ratio of luminosity to mass ( Houdek et al. 1999), and detections are more likely in magnetically inactive stars (Chaplin et al. 2011a), so the F-type stars in Table 1 should yield asteroseismic data comparable to Kepler targets in the magnitude range $Kp \sim 10$–11. The 27 day time series should also allow an independent check of the rotation periods determined from the Mount Wilson data.

We searched the García et al. (2014) rotation catalog for the fainter Kepler asteroseismic targets that can be considered analogs of the Mount Wilson stars in terms of both stellar properties and the expected data quality from TESS. The properties of KIC 10906929, the best Kepler analog of HD 212754, are listed in the fourth column of Table 2. KIC 10909629 has a $B-V$ color and rotation period that are both similar to HD 212754, but it is 5.25 mag fainter in the $V$ band. The chromospheric activity level of KIC 10909629 is not known, but the photospheric activity proxy $S_{\text{ph}}$ (Mathur et al. 2014) is comparable to that of KIC 10963065, which has $\log R'_{\text{HK}} < -5$. The gyro age from B07 (2.04 Gyr) and B10 (2.24 Gyr) are both significantly younger than the asteroseismic age (4.69 Gyr, Serenelli et al. 2017). The spectroscopic parameters and Gaia luminosity are similar to those of HD 212754, and support the conclusion that KIC 10909629 is substantially evolved despite its young gyro age. Again, this can be understood if the rotation period of KIC 10909629 stopped evolving after $\sim$2 Gyr. Based on these results, we predict that asteroseismic ages from TESS for the Mount Wilson stars in Table 1 will be significantly older than expected from gyrochronology.

The anticipated quality of the TESS observations for the Mount Wilson stars is illustrated in Figure 2. In the top panel, a 27 day segment of the long-cadence Kepler observations of KIC 10909629 clearly shows rotational modulation with a period near 12 days and a peak-to-peak amplitude of a few millimagnitudes. In the bottom panel, we show the power spectrum from 390 days of short-cadence data for KIC 10909629 (gray) and for the observations spanning 30 days (blue) that were originally used to detect solar-like oscillations in this star (Chaplin et al. 2011b). In both cases, the signatures of activity, granulation, and shot noise have been modeled and removed using the A2Z pipeline (Mathur et al. 2010). While the longer data set clearly shows the series of evenly spaced frequencies that are characteristic of solar-like oscillations, the shorter time series still reveals a significant power excess $\sim$900 $\mu$Hz and a signature of the regular spacing. When combined with spectroscopic parameters, these global properties of the oscillations are sufficient to constrain the stellar age with $\sim$10%–20% precision (Chaplin et al. 2014; Serenelli et al. 2017).

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Table 2

| KIC 9139151 | KIC 12009504 | KIC 10963065 | KIC 10909629 |
|----------------|----------------|----------------|----------------|
| $B-V$          | 0.520          | 0.556          | 0.509          | 0.540          |
| $p_{\text{rot}}$ [day] | 10.96 ± 2.22 | 9.39 ± 0.68 | 12.58 ± 1.70 | 12.37 ± 1.22 |
| $\log R'_{\text{HK}}$ | −4.954        | −4.977         | −5.054        | −5.054         |
| $t_{\text{gyro}}$ [Gyr] | 1.93 ± 0.37 | 1.06 ± 0.18 | 2.82 ± 0.57 | 2.04 ± 0.37 |
| $t_{\text{astero}}$ [Gyr] | 1.94 ± 0.31 | 3.44 ± 0.44 | 4.33 ± 0.30 | 4.69 ± 0.56 |
| $t_{\text{chro}}$ [Gyr] | 4.73          | 5.15           | 6.75           | ...           |
| $T_{\text{eff}}$ [K] | 6729          | 6179           | 6140           | 6265           |
| log $g$        | 4.38           | 4.21           | 4.29           | 3.90           |
| [$\text{Fe}/H]$ | +0.10          | −0.08          | −0.19          | −0.12          |
| $L/L_{\odot}$  | 1.71 ± 0.01    | 2.71 ± 0.01    | 1.93 ± 0.01    | 6.59 ± 0.11    |

References. García et al. (2014); Metcalfe et al. (2016); Creevey et al. (2017); Serenelli et al. (2017); Buchhave & Latham (2015); Gaia Collaboration et al. (2018).
gyrochronology. Using data directly from Barnes (2007, his Table 3), we demonstrate that the most significant differences between chromospheric ages and gyrochronology occur for the most evolved F-type dwarfs (Figure 1). We present several independent lines of evidence to corroborate this interpretation, including surface gravities, Gaia luminosities, and the predominant absence of activity cycles (Table 1). We identify analogs of these F-type stars among the sample of asteroseismic targets observed by Kepler, and we show that the asteroseismic ages agree with gyrochronology until a critical activity level \( \log R'_{\text{HK}} = -4.95 \) beyond which the two estimates diverge (Table 2). Finally, we use observations of a fainter analog from the Kepler sample to predict the quality of observations anticipated for these targets from the TESS mission, showing that rotation periods and solar-like oscillations should both be detectable (Figure 2). Considering our revised picture of the late stages of magnetic stellar evolution, we predict that these future observations will demonstrate that gyrochronology is unreliable for stars beyond the middle of their main-sequence lifetimes.

There are two key updates to the scenario for magnetic evolution outlined in this paper compared to that proposed by Metcalfe et al. (2016). First, it is now clear that the Rossby number from global convection simulations and that obtained from asteroseismic models that use a mixing-length prescription are not directly comparable (Brun et al. 2017). Based on solar determinations of the Rossby number from both methods, we now believe that the transition near \( \text{Ro} \sim 2 \) identified by van Saders et al. (2016) corresponds to a change in the
character of differential rotation seen near Ro \sim 1 in convection simulations. In the context of Metcalfe et al. (2016), this implies that stellar evolution across the Vaughan–Preston gap entirely precedes the magnetic transition, which occurs at a substantially lower activity level. Second, there is now observational evidence of possible anti-solar differential rotation in some stars that show higher than expected activity for their rotation rates (Brandenburg & Giampapa 2018). This suggests that when stars reach the critical Rossby number, the differential rotation might flip from solar-like to anti-solar (slow equator, fast poles). Conservation of angular momentum would require the shear to increase, leading to an enhancement of activity in the slowly rotating regime (Karak et al. 2015).

Future observations will determine whether this phenomenon represents a temporary phase, or an alternate pathway that coexists with the shutdown of the global dynamo.

While it may be disappointing that rotation is less useful as a diagnostic of age beyond the middle of stellar main-sequence lifetimes, the other Skumanich relation (activity–age) may not be similarly disrupted. Observations of chromospheric activity in a large sample of solar analogs suggest that, unlike rotation, the evolution of activity appears to be continuous across the magnetic transition (Lorenzo-Oliveira et al. 2018). Although the ages adopted for their analysis were derived from isochrones, the TESS mission is poised to provide reliable asteroseismic ages for bright stars down to V \sim 7 all around the sky. When combined with existing archives of chromospheric activity, it is possible that asteroseismic ages can be used to recalibrate the activity–age relation for older solar-type stars. Given the difficulty of obtaining time series measurements of the diminishing rotational modulation in such stars, and considering that minimal chromospheric variability makes one spectroscopic measurement more likely to be representative of the mean activity level, chromospheric activity might ultimately provide a more reliable age indicator for stars beyond middle-age.

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