LOOP GROUPS AND NONCOMMUTATIVE GEOMETRY

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Abstract. We describe the representation theory of loop groups in terms of K-theory and noncommutative geometry. This is done by constructing suitable spectral triples associated with the level $\ell$ projective unitary positive-energy representations of any given loop group $L_G$. The construction is based on certain supersymmetric conformal field theory models associated with $L_G$. We then generalize the construction, in the setting of conformal nets, to many other rational chiral conformal field theory models including coset models and the moonshine conformal net.

1. Introduction

Since its foundation by Connes [18], noncommutative geometry has been of growing importance, with impact on various fields of mathematics and physics: differential geometry, algebraic topology, index theory, quantum field theory, quantum Hall effect descriptions, etc., to name a few. The core idea is to work with an algebraic approach, which is in some sense “dual” to a topological one, namely instead of a given locally compact topological space one works with the commutative C*-algebra of continuous functions on it vanishing at infinity. It turns out that this offers several different and useful tools. According to Gelfand-Neimark’s theorem, every commutative C*-algebra is actually of this type. To make things meaningful, usually some additional structure is requested to be given alongside the C*-algebra of continuous functions. The step to the noncommutative setting consists now basically in still requiring this further structure but allowing for arbitrary noncommutative C*-algebras (or even more general noncommutative topological algebras) instead of the commutative C*-algebra of continuous functions (cf. [20, 46] for overview and a comprehensive study).

Examples of such further structures in noncommutative geometry are spectral triples $(A, (\pi, H), D)$, where apart from the algebra $A$, a representation $\pi$ on some Hilbert space $H$ and a self-adjoint operator $D$ with compact resolvent on $H$ are given. In the commutative case such a triple together with some additional data completely describes a smooth compact manifold, according to Connes’s recent reconstruction theorem [21]. In the noncommutative case, this is by far not enough to understand the complete structure well, but it suffices in order to partially understand the noncommutative geometry of given objects, and to compute K-homology classes, noncommutative Chern characters and index pairings with K-theory. There is a bivariant version of K-theory, called KK-theory: a bifunctor from the category of C*-algebras to abelian groups. KK-theory plays a fundamental role in the structure theory of C*-algebras and in noncommutative geometry; apart from the inherent group addition, it admits a so-called intersection product, and it generalizes both K-theory and K-homology. The operator algebraic nature of noncommutative geometry has enabled many fruitful connections to other areas in mathematics and physics (cf. again [20, 46]). In this article we establish a link with the representation theory of loop groups.

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Loop groups are well-studied examples of infinite-dimensional Lie groups \([71, 72]\). Given a smooth compact manifold \(X\) and a connected simply connected compact simple Lie group \(G\) with Lie algebra \(\mathfrak{g}\), a natural object to investigate is the group \(C^\infty(X,G)\) of smooth maps \(X \to G\), with point-wise multiplication and with the \(C^\infty\) topology (uniform convergence of all partial derivatives). It is an infinite-dimensional Lie group modeled on the topological vector space \(C^\infty(X,\mathfrak{g})\) of smooth maps \(X \to \mathfrak{g}\). Moreover, the Lie algebra of \(C^\infty(X,G)\) turns out to be \(C^\infty(X,\mathfrak{g})\) with point-wise brackets \([71, \text{Example 1.3}]\). Loop groups are obtained in the special case where \(X = S^1\). Accordingly the loop group of \(G\) is given by \(L G := C^\infty(S^1,G)\) and it is an infinite-dimensional Lie group with Lie algebra \(C^\infty(S^1,\mathfrak{g})\) \([72]\). The latter admits nontrivial central extensions corresponding to the affine Kac-Moody Lie algebras associated with \(\mathfrak{g}\) \([54, 72]\).

Loop groups are important objects to study in mathematics: they play a fundamental role in conformal quantum field theory and string theory. They have shown deep relations with various other mathematical areas such as number theory, subfactor theory, quantum groups and topological quantum field theories, see e.g. \([2, 31, 38, 41, 42, 85]\).

Loop groups also have a very interesting representation theory. It is obtained by restricting to a special class of (projective) unitary representations, namely the so-called positive-energy representations, or more precisely, those arising by integrating unitary highest weight irreducible representations of the underlying affine Kac-Moody algebra \(\hat{\mathfrak{g}}_C\). There is a distinguished central element \(c_\mathfrak{g}\) in \(\hat{\mathfrak{g}}_C\) and it takes positive integer value in such a representation; this integer is called the level of the representation. For every given level \(\ell\) there are only finitely many classes of such representations; remarkably, they generate a commutative ring \(R^\ell(L G)\), the so-called Verlinde fusion ring. The ring structure comes from the operator product expansion in the conformal field theory model associated to the representation theory of \(L G\) at level \(\ell\). Mathematically this can be described in various ways. For example the ring product can be defined through the so called Verlinde formula from the modular invariance property of the characters of the representations, see e.g. \([11, \text{Chap.5}]\) and \([12, \text{Chap. 6}]\). More conceptually \(R^\ell(L G)\) can be defined as the set of equivalence classes of objects in the braided tensor category of level \(\ell\) representations of \(L G\). The tensor structure on the latter category can be obtained e.g. through the Huang-Lepowsky theory of tensor products for vertex operator algebra modules \([52]\), see also \([72, \text{Chap. 6}]\). For explicit computations of the fusion rings see \([27]\).

In this paper we give a description of the Verlinde fusion ring \(R^\ell(L G)\) in terms of K-theory and noncommutative geometry for any given connected simply connected compact simple Lie group \(G\) and positive integer level \(\ell\). To this end we define a nonunital C*-algebra \(\mathcal{R}_{G,\ell}\) and a norm dense subalgebra \(\mathcal{R}_{G,\ell}^1 \subset \mathcal{R}_{G,\ell}\) associated with the level \(\ell\) positive-energy representations of \(L G\). Then, for any irreducible level \(\ell\) representation of \(L G\) we construct a spectral triple with algebra \(\mathcal{R}_{G,\ell}^1\) and show that the entire cyclic cohomology class of the corresponding JLO cocycle completely determines the unitary equivalence class of the loop group representation. These spectral triples give rise to a group isomorphism from the level \(\ell\) Verlinde fusion ring \(R^\ell(L G)\) of \(L G\) onto the K-homology group \(K^0(\mathcal{R}_{G,\ell})\) and, dually, to an isomorphism of \(R^\ell(L G)\) onto the K-theory group \(K_0(\mathcal{R}_{G,\ell})\). As a consequence we show that there is also an injective homomorphism of \(R^\ell(L G)\) into the KK-group \(KK(\mathcal{R}_{G,\ell}, \mathcal{R}_{G,\ell})\) such that the fusion product corresponds to the Kasparov product, i.e., a ring homomorphism. The construction of the above spectral triples is based on the supersymmetric chiral conformal field theory (CFT) models associated to the loop group representations. Although the above construction naturally identifies the additive group underlying the fusion ring \(R^\ell(L G)\) with the K-groups \(K^0(\mathcal{R}_{G,\ell})\) and \(K_0(\mathcal{R}_{G,\ell})\), in order to give the interpretation of the fusion product as a Kasparov product we need to know a priori the product in \(R^\ell(L G)\), which enters our construction as an external input. Accordingly, we have to
rely on one of the existing definitions of the product in $R^\ell(LG)$, e.g. the one coming from the theory of vertex operator algebras given in [52]. For this reason it is desirable to have an alternative definition of the fusion product which is more natural from an operator algebraic point of view. To this end, we give an interpretation of our construction in terms of conformal nets, which describes chiral CFTs in the operator algebraic approach to quantum field theory [49], in the light of the noncommutative geometrization program for CFT recently developed in [15, 13, 14, 16, 69], cf. also [74, 75] for related work. Actually, this program was one of the initial motivations for this work.

We obtain here for the first time examples for which this program can be carried out for all sectors of the corresponding conformal nets. This means that all irreducible sectors can be separated by JLO cocycles associated to spectral triples naturally arising from supersymmetry. By the results in Sect. 5.2 these examples include quite a large family of completely rational conformal nets. Moreover, the noncommutative geometric description of the representation theory of these conformal nets is directly related to the K-theoretic description recently given in [11, 12]. In order to interpret our construction this way we have to assume that the conformal net associated with the representation theory of $LG$ at level $\ell$ is completely rational and that the corresponding ring generated by the Doplicher-Haag-Roberts (DHR) endomorphisms is isomorphic to $R^\ell(LG)$. This assumption is known to be satisfied e.g. for $G = SU(n)$ at any positive integer level, and it is widely expected to be true in general although this is still a very important open problem which goes far beyond the scope of this paper, see e.g. [57, Problem 3.32]. If the pair $(G, \ell)$ satisfies this assumption then the ring isomorphism from $R^\ell(LG)$ into $KK(\mathbb{K}_G^\ell, \mathbb{K}_G^\ell)$ arises naturally from the fact that DHR endomorphisms of the net give rise to endomorphisms of the C*-algebra $\mathbb{K}_G^\ell$ and hence to KK-classes in $KK(\mathbb{K}_G^\ell, \mathbb{K}_G^\ell)$, [11, 12].

A different K-theoretical description of the fusion ring $R^\ell(LG)$ has been given by Freed, Hopkins and Teleman (FHT) [34, 35, 36, 37]. In those papers the Verlinde fusion ring of $LG$ at level $\ell$ has been identified with a twisted version of the equivariant topological K-theory group $K_G(LG)$, where $G$ acts on itself by conjugation. The twisting is determined by the level $\ell$. Under this identification the fusion product corresponds to the convolution (Pontryagin) product on the twisted $K_G(LG)$ and the latter can be defined without direct reference to the Verlinde fusion product in $R^\ell(LG)$ [36, Thm.1], see also [35]. Very interesting relations of this result with subfactor theory and modular invariants have been investigated by Evans and Gannon [28, 29, 30].

Our results indirectly show the identification of the above twisted equivariant K-theory of $G$ with the K-theory group $K^0(\mathbb{K}_G^\ell)$ of the C*-algebra $\mathbb{K}_G^\ell$. Moreover there are various structural similarities between our approach and FHT. Our Dirac operators are essentially the same as those considered in [36, Sec. 11], see also [37]. In our case the Dirac operators are used to construct spectral triples with algebra $\mathbb{K}_G^\ell$ in order to obtain K-homology classes $K^0(\mathbb{K}_G^\ell)$. In the FHT case the Dirac operators are used to construct equivariant Dirac families which give rise to twisted equivariant K-classes.

We believe it would be very interesting to exploit this relationship deeper and more explicitly in future. In particular we believe that a better understanding of this relationship would shed more light on these K-theoretical approaches to the representation theory of loop groups and in particular to its deep mathematical relations with conformal field theory. More generally, it could reveal new connections between twisted K-theory, non-commutative geometry and subfactor theory. Moreover it could shed some light on the relation of the FHT convolution product in equivariant twisted K-theory and the Kasparov product defined in [80]. From this point of view our analysis should not be seen as alternative to the FHT work but rather complementary to it. It should be pointed out that the FHT analysis goes beyond the case in which $G$ is
a simply connected compact Lie groups and actually it also deals with non-connected compact Lie groups (in particular with finite groups) and twisted loop groups. On the other hand we give in this paper an abstract version of our construction, formulated in terms of rational conformal nets, which applies to many conformal field theory models besides the case of loop groups. In particular we can cover the case of loop groups associated to arbitrary connected compact Lie groups but also coset models, minimal models and the conformal net analogue of the Frenkel-Lepowsky-Meurman moonshine vertex operator algebra \cite{39} constructed by Kawahigashi and Longo in \cite{60}. We should also point out that the K-theoretical description of the DHR fusion ring in \cite{11, 12} is very general since it works for all completely rational conformal nets.

Our paper is organized as follows. We start off with a preliminary section on noncommutative geometry for nonunital algebras, which is difficult to find in literature in this form.

The main Section 3 then deals with the construction of our noncommutative geometric objects: for a given connected simply connected compact simple Lie group $G$ and hence for its corresponding loop group $L G$, we construct Dirac operators and a global differentiable algebra $\mathcal{R}_{G_{\ell}}$ on which the Dirac operators act; they give rise to spectral triples, noncommutative Chern characters (also known as JLO cocycles) and index parings with K-theory classes corresponding to characteristic projections in $\mathcal{R}_{G_{\ell}}$. The underlying idea in the construction is supersymmetry: tensoring the loop group by the (graded) CAR algebra, we turn given representations into graded representations; for the latter special representations, the so-called super-Sugawara (or Kac-Todorov) construction \cite{56} guarantees supersymmetry and Dirac operators. The differentiable algebra $\mathcal{R}_{G_{\ell}}$ comes out as a natural byproduct. Moreover, the universal coefficient theorem then identifies the fusion product as a Kasparov product at an abstract level.

In Section 4, we discuss the CFT interpretation of the results in Section 3 in terms of conformal nets. After a quick introduction to conformal nets, we show that under natural assumptions, for $G$ simply connected, the ring isomorphism of $\mathcal{R}^\ell(L G)$ into $KK(\mathcal{R}_{G_{\ell}}, \mathcal{R}_{G_{\ell}})$ can be interpreted in terms of DHR endomorphisms, a fact that merges nicely with the results in \cite{11, 12, 14, 15}.

In Section 5, we introduce the notion of superconformal tensor product for conformal nets that can be considered as a generalization of the super-Sugawara construction for loop groups models. We show that our previous analysis for loop groups $L G$ with $G$ simply connected, extends in a rather straightforward way to the case of completely rational conformal nets with superconformal tensor product. We then provide many examples of completely rational conformal nets admitting a superconformal tensor product as announced above.

2. Entire cyclic cohomology for nonunital Banach algebras

Let $(A, \|\cdot\|)$ be a Banach algebra. If $A$ is nonunital, its unitization $\hat{A}$ is obtained by adjoining a unit $1_{\hat{A}}$ to $A$, with multiplication defined by
\[(a_1 + t_1 1_{\hat{A}}) \cdot (a_2 + t_2 1_{\hat{A}}) := a_1 a_2 + t_1 a_2 + t_2 a_1 + t_1 t_2 1_{\hat{A}},\]
and norm $\|a + t 1_{\hat{A}}\| := \|a\| + |t|$, for every $a + t 1_{\hat{A}} \in \hat{A}$. Then $A \subset \hat{A}$ is a closed ideal and $\hat{A}/A \simeq \mathbb{C}$. If $\|\cdot\|$ was a C*-norm, one can construct a C*-norm on $\hat{A}$ equivalent to $\|\cdot\|$ \cite{41 II.1.2}. However, we will not need this. Moreover, we shall drop the superscript “$\sim$” of $\|\cdot\|\sim$ henceforth; analogously for subscripts of $\|\cdot\|$ indicating the algebra whenever they appear and confusion is unlikely. If $\pi : A \to B(\mathcal{H}_\pi)$ is a representation, then it extends to a unital representation $\tilde{\pi}$ of $\hat{A}$ by $\tilde{\pi}(a + t 1_{\hat{A}}) := \pi(a) + t 1_{\mathcal{H}_\pi}$ and we shall always work with this extension.

For every $r \in \mathbb{N}$, let $M_r(A)$ be the algebra of $r \times r$ matrices with entries in $A$. The maps $x \mapsto \text{diag}(x, 0)$ define natural embeddings $M_r(A) \to M_{r+1}(A), r \in \mathbb{N}$. Then
$M_\infty(A)$ is defined as the corresponding inductive limit, i.e., the algebra of infinite-dimensional matrices with entries in $A$ and all but finitely many of them zero [3 5.1]. The algebras $M_r(A), r \in \mathbb{N}$, can be made into Banach algebras in many equivalent ways. One can choose the norms on $M_r(A)$ in such a way that the embeddings are isometries. Here we fix any sequence of such norms. It induces a norm on $M_\infty(A)$ which is accordingly made into a normed algebra, too. What follows in this paper does not depend on the particular choice of the norms.

$A$ is called stably-unital if $M_\infty(A)$ has an approximate identity of idempotents; this is e.g. the case if $A$ has an approximate identity of idempotents [3, 5.5].

Entire cyclic cohomology is usually studied in the setting of unital Banach algebras. A definition for nonunital ones using the cyclic cocomplex can be found in detail in [22 Sect.6]. The subsequent one following [65 Sect.2] is stated in the $(b, B)$-bicomplex, which is more suitable in our setting since we want to apply it to the JLO cochain. The two definitions are equivalent according to [65 Sect.2] together with [67 Prop.4.2].

**Definition 2.1.** (i) Let $A$ be a nonunital Banach algebra and, for any nonnegative integer $n$, let $C^n(A)$ be the vector space of reduced $(n+1)$-linear forms $\phi_n$ on the unitalization $\tilde{A}$, i.e., such that

- $\phi_0(1_{\tilde{A}}) = 0$,
- $\phi_n(a_0, a_1, \ldots, a_n) = 0$ if $a_i = 1_{\tilde{A}}$ for some $1 \leq i \leq n$ (simplicial normalization [8]).

For integers $n < 0$ we set $C^n(A) := \{0\}$. Let $C^\bullet(A) := \prod_{n=0}^{\infty} C^n(A)$ be the space of sequences $\phi = (\phi_n)_{n \in \mathbb{N}_0}$, with $\phi_n \in C^n(A)$ and define the operators $b : C^\bullet(A) \to C^\bullet(A)$ and $B : C^\bullet(A) \to C^\bullet(A)$ by

$$(b\phi)_n(a_0, \ldots, a_n) := \sum_{j=0}^{n-1} (-1)^j \phi_{n-1}(a_0, \ldots, a_j a_{j+1}, \ldots, a_n) + (-1)^n \phi_{n-1}(a_n a_0, a_1, \ldots, a_{n-1}),$$

$$(B\phi)_n(a_0, \ldots, a_n) := \sum_{j=0}^{n} (-1)^j \phi_{n+1}(1_{\tilde{A}}, a_j, \ldots, a_n, a_0, \ldots, a_{j-1}), \quad a_0, \ldots, a_n \in \tilde{A}.$$

The linear map $\partial : C^\bullet(A) \to C^\bullet(A)$ defined by $\partial := b + B$ satisfies $\partial^2 = 0$ and, with the coboundary operator $\partial$, $C^\bullet(A)$ becomes the cyclic cocomplex $C^\bullet(A) = (C^c(A), C^o(A))$ over $\mathbb{Z}/2\mathbb{Z}$, namely the elements of $C^c(A) = \prod_{n=0}^{\infty} C^{2n}(A)$ (the even cochains) are mapped into the elements of $C^o(A) = \prod_{n=0}^{\infty} C^{2n+1}(A)$ (the odd cochains) and vice versa. The elements $\phi \in C^\bullet(A)$ satisfying $\partial \phi = 0$ are called the cyclic cocycles of $A$.

(ii) A cochain $\phi = (\phi_n)_{n \in \mathbb{N}_0} \in C^\bullet(A)$ is called entire if

$$|\phi_n(a_0, \ldots, a_n)| \leq \frac{1}{\sqrt{n!}} \|a_0\| \cdots \|a_n\|, \quad a_i \in \tilde{A}, n \in \mathbb{N}_0.$$

Letting $CE^\bullet(A)$ be the entire elements in $C^\bullet(A)$, one defines the entire cyclic cohomology $(HE^c(A), HE^o(A))$ of $A$ as the cohomology of the cocomplex $((CE^c(A), CE^o(A)), \partial)$. The cohomology class of an entire cyclic cocycle $\phi \in CE^\bullet(A) \cap \ker(\partial)$ will be denoted by $[\phi]$.

We recall that in the case of unital $A$ instead no unitalization is considered; $C^n(A)$ stands then for the simplicially normalized (not necessarily reduced) cochains on $A$ itself, and entire cyclic cohomology is defined accordingly [22].

**Definition 2.2.** [3 Sect.5.5] Let $A$ be a stably-unital Banach algebra. We denote by $PM_r(A)$ the set of idempotents in $M_r(A), r \in \mathbb{N} \cup \{\infty\}$. An equivalence relation on
$P M_\infty(A)$ is defined by $p \sim q$ if there are $x, y \in M_\infty(A)$ such that $p = xy$, $q = yx$. There is a binary operation

$$(p_1, p_2) \in P M_{r_1}(A) \times P M_{r_2}(A) \mapsto p_1 \oplus p_2 := \text{diag}(p_1, p_2) \in P M_{r_1+r_2}(A),$$

which turns $P M_\infty(A)/\sim$ into an abelian semigroup. Then the $K_0$-group of $A$ is defined as

$$K_0(A) := \text{Grothendieck group of } P M_\infty(A)/\sim.$$  

We write $[p]$ for the element in $K_0(A)$ induced by a projection $p \in A$.

**Definition 2.3.** A $\theta$-summable even spectral triple is a triple $(A, (\pi, \mathcal{H}), D)$, where

- $A$ is an algebra;
- $\mathcal{H}$ is a Hilbert space graded by a selfadjoint unitary $\Gamma$, and $\pi$ is a representation of $A$ on $\mathcal{H}$ commuting with $\Gamma$;
- $D$ is a self-adjoint operator with $e^{-tD^2}$ trace-class for all $t > 0$, with $\Gamma D \Gamma = -D$, and with $\pi(A) \subset \text{dom}(\delta_D)$, where $\delta_D$ is the derivation on $B(\mathcal{H})$ induced by $D$ as explained below.

Note that, according to the above definition, if $A$ is nonunital then $(A, (\pi, \mathcal{H}), D)$ is a $\theta$-summable even spectral triple if and only if $(\tilde{A}, (\tilde{\pi}, \mathcal{H}), D)$ is. Henceforth we shall deal with spectral triples in the case where $A$ is a nonunital Banach algebra.

Given a self-adjoint operator $D$ on $\mathcal{H}$, one associates a derivation $\delta_D$ of $B(\mathcal{H})$ as follows: define $\text{dom}(\delta_D)$ as the algebra of $x \in B(\mathcal{H})$ such that

$$xD \subset Dx - y$$

for some $y \in B(\mathcal{H})$, and $\delta_D(x) := y$ in this case, cf. [7, Sect.3.2].

The JLO cochain $\tau = (\tau_n)_{n \in \mathbb{Z}_n}$ of a spectral triple $(A, (\pi, \mathcal{H}), D)$ with $A$ a nonunital Banach algebra is defined as $\tau := \tilde{\tau} - \psi$, where

$$\tilde{\tau}_n(a_0, \ldots, a_n) = \int_{0 \leq t_1 \leq \ldots \leq t_n \leq 1} \text{tr} \left( \Gamma \tilde{\pi}(a_0) e^{-t_1 D^2} \delta_D(\tilde{\pi}(a_1)) e^{-t_2 D^2} \delta_D(\tilde{\pi}(a_2)) \cdots \right)$$

$$\cdots \delta_D(\tilde{\pi}(a_{n-1})) e^{-t_n D^2} \right) dt_1 \cdots dt_n$$

and

$$\psi_n(a_0, \ldots, a_n) := \begin{cases} t\tilde{\pi}_0(1_{\tilde{A}}) & \text{if } n = 0, \ a_0 \in t1_{\tilde{A}} + A \\ 0 & \text{if } n > 0, \end{cases}$$

for every $a_i \in \tilde{A}$ and every $n \in \mathbb{N}_0$. Clearly, in restriction to entries in $A$, we have $\tau \restriction_A = \tilde{\tau} \restriction_A$.

Notice that in the common context of unital algebras and spectral triples (e.g. in [15]), $\tilde{\tau}$ would be a simplicially normalized (but not reduced) cochain on $A$ itself, namely the original JLO cochain from [53].

Now, let $(A, (\pi, \mathcal{H}), D)$ be a $\theta$-summable even spectral triple with grading operator $\Gamma$. Let $\mathcal{H}_\pm = \ker(\Gamma \mp 1_{B(\mathcal{H})})$ so that we have the decomposition $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-$. If $T$ is a densely defined operator on $\mathcal{H}$ which is odd, i.e., such that $\Gamma T \Gamma = -T$, then we can write

$$T = \begin{pmatrix} 0 & T_- \\ T_+ & 0 \end{pmatrix},$$

with operators $T_\pm$ from (a dense subspace of) $\mathcal{H}_\pm$ to $\mathcal{H}_\mp$. Accordingly, if $(A, (\pi, \mathcal{H}), D)$ is a $\theta$-summable even spectral triple then for the selfadjoint $D$ we can write

$$D = \begin{pmatrix} 0 & D_- \\ D_+ & 0 \end{pmatrix},$$
with $D_- = D_+^*$. On the other hand, if $T$ is even, i.e., it commutes with $\Gamma$, then we can write

$$T = \begin{pmatrix} T_+ & 0 \\ 0 & T_- \end{pmatrix},$$

with operators $T_\pm$ from (a dense subspace of) $\mathcal{H}_\pm$ to $\mathcal{H}_\pm$.

For any positive integer $r$ we denote by $\sigma_r$ the representation of $M_r(A)$ on $\mathcal{H}_r := \mathbb{C}^r \otimes \mathcal{H}$ defined by $\sigma_r(m \otimes a) := m \otimes \pi(a)$, $m \in M_r(\mathbb{C})$, $a \in A$. Moreover, for every operator $T$ on $\mathcal{H}$ we consider the operator $T_r := 1 \otimes T$ on $\mathcal{H}_r$. Then, for every $r \in \mathbb{N}$, $(M_r(A), (\sigma_r, \mathcal{H}_r), D_r)$ is a $\theta$-summable even spectral triple which is even with grading operator $\Gamma_r$. For an $n$-linear form $\phi_{n-1} : \hat{A}^\otimes n \to \mathbb{C}$, we introduce the $n$-linear form $\phi_{n-1}^r$ on $(M_r(\mathbb{C}) \otimes \hat{A})^\otimes n = M_r(\hat{A})^\otimes n$ as

$$\phi_{n-1}^r(m_0 \otimes a_0, \ldots, m_n \otimes a_n) := \text{tr}(m_0 \ldots m_n)\phi_{n-1}(a_0, \ldots, a_n)$$

and linear extension.

The following fundamental theorem is an adaptation to the present nonunital setting of the corresponding unital setting, cf. e.g. [19, 53, 43, 15].

**Theorem 2.4.** Let $(A, (\pi, \mathcal{H}), D)$ be a $\theta$-summable even spectral triple such that $A$ is a nonunital Banach algebra and the representation $\pi$ of $A$ in the Banach algebra $(\text{dom}(\delta_D), \| \cdot \| + \| \delta_D(\cdot)\|_{B(\mathcal{H})})$ is continuous.

(i) The cochain $(\tau_r)_{n \in 2\mathbb{N}_0}$ is an even entire cyclic cocycle on $A$. We call it the JLO cocycle or Chern character of $(A, (\pi, \mathcal{H}), D)$.

(ii) The values of the maps $(\phi, p) \in (CE^e(A) \cap \ker(\partial)) \times PM_r(A) \mapsto \phi(p) \in \mathbb{C}$, $r \in \mathbb{N}$, where

$$\phi(p) := \phi^r_0(p) + \sum_{k=1}^{\infty} (-1)^k \frac{(2k)!}{k!} \phi^r_{2k}((p - \frac{1}{2}), p, \ldots, p),$$

only depend on the cohomology class $[\phi]$ of the entire cyclic cocycle $\phi$ and on the $K_0(\hat{A})$-class $[p]$ of the idempotent $p$ and are additive on the latter. If $A$ is stably-unital, this gives rise to a pairing $\langle [\phi], [p] \rangle := \phi(p)$ between the even entire cyclic cohomology $HE^e(A)$ and K-theory $K_0(A)$. Moreover, the operator $\pi_r(p)_-D_r + \pi_r(p)_+$ from $\pi_r(p)_-\mathcal{H}_r+ \to \pi_r(p)_-\mathcal{H}_r+$ is a Fredholm operator and for the even JLO cocycle $\tau$ we have

$$\tau(p) = \langle [\tau], [p] \rangle = \text{ind}_{\pi_r(p)_-\mathcal{H}_r+} (\pi_r(p)_-D_r + \pi_r(p)_+) \in \mathbb{Z}.$$

**Proof.** For unital algebras and spectral triples this theorem is well-known, namely the JLO cochain $\hat{\tau}$ is an even entire cyclic cocycle on the unital algebra $\hat{A}$ [20, 43, 53]. The same is true for $\psi$ as is easily verified. $\tau = \hat{\tau} - \psi$ is again an entire cyclic cocycle of $\hat{A}$ because both $\hat{\tau}$ and $\psi$ are so. Evenness is obvious, i.e., $\tau_n = 0$ if $n \in \mathbb{N}_0$ is odd. For every $n \in \mathbb{N}_0$, $\tau_n$ is reduced, i.e., $\tau_0(1_{\hat{A}}) = 0$ and $\tau_n(a_0, a_1, \ldots, a_n) = 0$ if $a_i = 1_{\hat{A}}$ for some $1 \leq i \leq n$. Thus $\tau$ is an even entire cyclic cocycle of $A$, proving (i).

Part (ii) becomes clear, too, by considering first all the statements and the well-known pairing between entire cyclic cohomology and K-theory (cf. e.g. [19, 20, 53, 43, 15]) for the unital algebra $\hat{A}$, namely: $\tau$ pairs with $K_0(\hat{A})$ as

$$\tau(p) := \tau_0^r(p) + \sum_{k=1}^{\infty} (-1)^k \frac{(2k)!}{k!} \tau_{2k}^r((p - \frac{1}{2}), p, \ldots, p),$$

if $p \in M_r(\hat{A})$ represents a class in $K_0(\hat{A})$, and $\tau(p_1) = \tau(p_2)$ for projections $p_1, p_2 \in M_\infty(\hat{A})$ if $p_1$ and $p_2$ belong to the same class in $K_0(\hat{A})$. Finally, if $A$ is stably-unital, then $K_0(A)$ is defined according to Definition 2.2 using equivalence classes of projections in $M_\infty(A)$; if moreover $p_1, p_2 \in M_\infty(A) \subset M_\infty(\hat{A})$ and give rise to the same class in $K_0(\hat{A})$ then according to [34, 5.5.2] they have the same class in $K_0(\hat{A})$, so $\tau(p_1) = \tau(p_2)$.
according to the previous step. Thus the pairing of \( \tau \) with \( K_0(A) \) restricts to a pairing with \( K_0(A) \). The Fredholm property and the final formula follow directly from the unital case by restriction.

3. JLO COCYCLE AND AN INDEX PAIRING FOR LOOP GROUPS

Let \( G \) be a connected simply connected compact simple Lie group and denote its Lie algebra by \( \mathfrak{g} \) and its dimension by \( d \in \mathbb{N} \), and let \( L G := C^\infty(S^1, G) \) the corresponding loop group. It is an infinite-dimensional Lie group, with Lie algebra the loop algebra \( L \mathfrak{g} := C^\infty(S^1, \mathfrak{g}) \). We denote by \( L \mathfrak{g}_C \) the complexification of \( L \mathfrak{g} \). The Lie subalgebra \( \mathfrak{g}_C \subset L \mathfrak{g}_C \) consisting of elements with finite Fourier series admits a nontrivial central extension \( \mathfrak{g}_C \) called the affine Kac-Moody algebra associated with \( \mathfrak{g} \), see \([41, 54, 72]\). We shall denote the corresponding distinguished central element by \( c_\mathfrak{g} \).

We would like to consider positive-energy representations of \( L G \) from the point of view of noncommutative geometry. A strongly continuous projective unitary representation \( \lambda : L G \to U(H_\lambda)/U(1) \) on a Hilbert space \( H_\lambda \) is of positive energy if there is a strongly continuous one-parameter group \( U^\lambda : \mathbb{R} \to U(H_\lambda) \) whose self-adjoint generator \( L^\lambda_0 \) (the conformal Hamiltonian) has non-negative spectrum and such that

\[
U^\lambda(t)\lambda(g)U^\lambda(t)^* = \lambda(g_t), \quad g \in L G,
\]

where \( g_t \) is defined by \( g_t(z) := ge^{-it}z \), \( z \in S^1 \). Eq. (3.1) should be understood in the projective sense so that if for every \( g \in L G \), \( \lambda(g) \in U(H_\lambda) \) denotes a given choice of a representative of \( \lambda(g) \in U(H_\lambda)/U(1) \) then

\[
U^\lambda(t)\hat{\lambda}(g)U^\lambda(t)^* = \chi(g, t)\hat{\lambda}(g_t), \quad g \in L G,
\]

with \( \chi(g, t) \in U(1) \). Note that

\[
\chi(g, 2\pi)1_{H_\lambda} = U^\lambda(2\pi)\hat{\lambda}(g)U^\lambda(2\pi)^*\hat{\lambda}(g)^* = \lambda^\vee
\]

so that \( g \mapsto \chi(g, 2\pi) \) is a continuous character of \( L G \) and hence \( \chi(g, 2\pi) = 1 \) for all \( g \in L G \) because \( L G \) is a perfect group \([72, 3.4.1]\). Hence, if \( \lambda \) is irreducible, then \( U^\lambda(\pi) = e^{2\pi i L^\lambda_0} \) is a multiple of the identity and hence \( t \mapsto U^\lambda(t) \) factors to a representation of \( Rot(S^1) \), the group of rotations of \( S^1 \).

If \( \pi \) is an irreducible unitary highest weight representation of \( \mathfrak{g}_C \), cf. \([51]\), then it exponentiates to an irreducible projective unitary positive-energy representation \( \lambda_\pi \) of \( L G \) \([45, Thm.6.7], [78, Thm.6.1.2]\). Moreover, there is a canonical choice of the positive operator \( L^\lambda_0^\vee \). The value \( \ell := \pi(c_\mathfrak{g}) \in \mathbb{N} \) is called the level of the representation \( \lambda_\pi \). In this paper we shall always deal with direct sums of such representations \( \lambda_\pi \) at fixed level. It has been shown by A. Wassermann that these are exactly the irreducible strongly continuous projective unitary positive-energy representations \( \lambda \) of \( L G \) such that the unitary one-parameter group \( U^\lambda \) implementing the rotations action on \( L G \) is diagonalizable with finite-dimensional eigenspaces, see \([79, Thm. 4.1]\). Equivalently they are the irreducible smooth strongly continuous projective unitary positive-energy representations of \( L G \) \([72, Sec. 9.3]\). Any of these representations determines a central extension of \( L G \) by \( U(1) \) \([72]\) and the equivalence class of the corresponding circle bundle on \( L G \) only depends on the level \([72]\). Accordingly, the representations of fixed level \( \ell \) correspond to the representations determining the associated central extension.

The Verlinde fusion ring \( R^\ell(L G) \) is the ring of formal differences (Grothendieck ring) associated with the semiring of equivalence classes of finite direct sums of such representations \( \lambda_\pi \) at level \( \ell \) \([31, 35, 36, 37, 27]\). The product is the so called fusion product. For any \( \ell \), there are only finitely many (say \( N \)) equivalence classes of irreducible unitary highest weight level \( \ell \) representations of \( \mathfrak{g}_C \). They are all described
and the corresponding loop group representations are discussed in \[72\] Sect.9, cf. also \[78\] Sect.6.

Given now a fixed level \(\ell\), let \((\lambda_i, \mathcal{H}_\lambda)\), with \(i = 0, \ldots, N - 1\), denote an arbitrary fixed maximal family of mutually inequivalent irreducible representations as above at level \(\ell\), with \(\lambda_0\) the vacuum representation, i.e., the representation corresponding to integral highest weight 0. Then every element in \(R^\ell(LG)\) can be written as

\[
\sum_{i=0}^{N-1} m_i[\lambda_i], \quad m_i \in \mathbb{Z}.
\]

The fusion product is determined by the Verlinde fusion rule coefficients \(N^k_{ij} \in \mathbb{N}_0\) such that

\[
[\lambda_i] \cdot [\lambda_j] = \sum_{k=0}^{N-1} N^k_{ij} [\lambda_k],
\]

for all \(i, j = 0, \ldots, N - 1\).

For every \(i = 0, \ldots, N - 1\), one can consider the conjugate representation \(\lambda_i^c\) of \(\lambda_i\), which is uniquely determined by the condition

\[
N^0_{ij} = \delta_{i,j}, \quad j = 0, \ldots, N - 1.
\]

The fusion rule coefficients also satisfy the identities

\[
N^k_{ij} = N^k_{ji},
\]

and

\[
N^k_{ij} = N^k_{ij} = N^k_{ij}.
\]

Now, let \((\lambda_{\text{red}}, \mathcal{H}_{\text{red}})\) denote the direct sum of all \((\lambda_i, \mathcal{H}_\lambda)\). By construction, every representation in consideration is unitary equivalent to a subrepresentation of a suitable multiple of \(\lambda_{\text{red}}\). We call \(\lambda_{\text{red}}\) the reduced universal representation or simply reduced representation. We now define two C*-algebras associated with the level \(\ell\) representations of \(LG\): the reduced universal algebra

\[
B_{\ell} := \lambda_{\text{red}}(LG)^\ddagger \simeq \bigoplus_{i=0}^{N-1} B(\mathcal{H}_\lambda_i)
\]

and the compact universal algebra

\[
\mathcal{R}_{\ell} := B_{\ell} \cap K(\mathcal{H}_{\text{red}}).
\]

We write \(\pi_\lambda\) for the (unique normal) representation of \(B_{\ell}\) such that \(\lambda_i = \pi_\lambda \circ \lambda_{\text{red}}\), and continue to use the same symbol for its restriction to \(\mathcal{R}_{\ell}\).

Given the Hilbert space \(\mathcal{K} := L^2(S^1, \mathbb{C}^{2d})\) with complex conjugation operator \(\gamma\), the corresponding self-dual CAR algebra \(\text{CAR}(\mathcal{K}, \gamma)\) (\[1\], sometimes also called algebra of 2d free fermionic fields) is the unital graded C*-algebra generated by odd \(F(f)\), for \(f \in \mathcal{K}\), such that \([F(f), F(g)]_+ = \langle f, g \rangle \mathbb{1}\) and \(F(f)^* = F(\gamma f) = F(\bar{f})\). It has so-called Ramond and Neveu-Schwarz irreducible representations, and we are interested in the former one \((\pi_R, \mathcal{H}_{\pi_R})\), which is graded by a selfadjoint unitary \(\Gamma_{\pi_R}\) because the dimension \(2d\) is even. The rotation group acts naturally on \(\text{CAR}(\mathcal{K})\), and its infinitesimal generator (also called conformal Hamiltonian) \(L^R_{0\ell}\) in the representation \(\pi_R\) has positive discrete spectrum. The smallest eigenvalue of \(L^R_{0\ell}\) is given by \(h_R := d/8\), and the corresponding eigenspace has dimension \(2^d\). We write \(\mathcal{H}_{\pi_R, 0, +}\) for the even part of this eigenspace with respect to \(\Gamma_{\pi_R}\), which has dimension \(2^{d-1}\). Let henceforth \(e_R\) be the projection onto an arbitrary but fixed one-dimensional subspace of \(\mathcal{H}_{\pi_R, 0, +}\). For a more expanded summary about the CAR algebra and Ramond representations with the present notation, we refer to \[15\] Sect.6 and for details and proofs to \[1\] \[6\] together with \[96\] Sect.5.3.
For any \( i \in \{0, \ldots, N-1\} \) we fix a lowest energy unit vector \( \Omega_{\lambda_i} \in \mathcal{H}_{\lambda_i} \) and denote by \( q_{\Omega_i} \) the orthogonal projection onto \( \mathbb{C}\Omega_{\lambda_i} \). Now, for every representation \((\lambda_i; \mathcal{H}_{\lambda_i})\) of \( LG \), we define the degenerate representation \( \hat{\pi}_{\lambda_i} : B_{G_e} \to B(\hat{\mathcal{H}}_{\lambda_i}) \) on the Hilbert space \( \mathcal{H}_{\lambda_i} := \mathcal{H}_{\lambda_i} \otimes \mathcal{H}_{\lambda_0} \otimes \mathcal{H}_{\pi_R} \) by
\[
\hat{\pi}_{\lambda_i}(x) := \pi_{\lambda_i}(x) \otimes q_{\Omega_{\lambda_0}} \otimes e_R, \quad x \in B_{G_e}.
\] (3.8)
This way, \( \hat{\mathcal{H}}_{\lambda_i} \) is graded by a grading operator \( \hat{\Gamma}_{\lambda_i} := 1_{\lambda_i} \otimes 1_{\lambda_0} \otimes \pi_{\Gamma_R} \) and \( \hat{\pi}_{\lambda_i} \) is even, i.e., commutes with \( \text{Ad} \hat{\Gamma}_{\lambda_i} \). On \( \mathcal{H}_{\lambda_i} \), we can define the total conformal Hamiltonian
\[
\hat{L}^{\lambda_i} := L_0^{\lambda_i} \otimes 1_{\mathcal{H}_{\lambda_0}} \otimes 1_{\mathcal{H}_{\pi_R}} + 1_{\mathcal{H}_{\lambda_i}} \otimes L_0^{\lambda_0} \otimes 1_{\mathcal{H}_{\pi_R}} + 1_{\mathcal{H}_{\lambda_i}} \otimes 1_{\mathcal{H}_{\lambda_0}} \otimes L_0^{\pi_R}.
\] (3.9)
It satisfies
\[
e^{iL_0^{\lambda_i}} \left( \pi_{\lambda_i}(x) \otimes q_{\Omega_{\lambda_0}} \otimes e_R \right) e^{-iL_0^{\lambda_i}} = \left( e^{iL_0^{\lambda_i}} \pi_{\lambda_i}(x) e^{-iL_0^{\lambda_i}} \right) \otimes q_{\Omega_{\lambda_0}} \otimes e_R
\] so that it generates the rotation action in the representation \( \hat{\pi}_{\lambda_i} \).

The following proposition shows that \( \hat{\pi}_{\lambda_i} \) can be considered as a “supersymmetric representation”, cf. [15][16].

**Proposition 3.1.** For \( i = 0, \ldots, N-1 \), there is an odd selfadjoint operator \( D_{\lambda_i} \) on \( \mathcal{H}_{\lambda_i} \) such that \( D_{\lambda_i}^2 = \hat{L}^{\lambda_i} - \frac{c}{24} \hat{\eta}_{\lambda_i} \), with \( c = 2(\frac{d}{2} + \frac{df}{e+h}) \), where \( h^\vee \) is the dual Coxeter number of \( g_C \). Furthermore, the spectrum of \( D_{\lambda_i} \) does not contain 0.

**Proof.** For the proof we will need a couple of background facts. For complete details we refer to [78] together with [72, Ch.9] and [54, 56]. The notation used here has been introduced and explained in [15, Sect.6].

By construction, the positive-energy representation \( \lambda_i \) of \( LG \) at level \( \ell \) comes from a representation of the affine Kac-Moody algebra \( \hat{g}_C \) on (the dense subspace of finite energy vectors of) \( \mathcal{H}_{\lambda_i} \) by integration, denoted again by \( \lambda_i \). Let us consider the representation
\[
\Lambda_i := \lambda_i \otimes \lambda_0 \otimes \pi_R : \hat{g}_C \oplus \hat{g}_C \oplus \text{CAR}(\mathcal{K}, \gamma) \to B(\hat{\mathcal{H}}_{\lambda_i}).
\]
Its generators are given by even \( J_{a}^{\Lambda_i} := J_a^{\Lambda_i} \otimes 1_{\lambda_0} \otimes 1_{\pi_R} \) for \( a = 1, \ldots, d \), \( J_{a-d}^{\Lambda_i} := 1_{\lambda_i} \otimes J_{a-d} \otimes 1_{\pi_R} \) for \( a = d+1, \ldots, 2d \), and odd \( F_{a}^{\Lambda_i} := 1_{\lambda_i} \otimes 1_{\lambda_0} \otimes F_{a}^{\pi_R} \), with \( a = 1, \ldots, 2d \) and \( n \in \mathbb{Z} \) (we shall henceforth drop the “\( \otimes 1 \)” if confusion is unlikely), and they satisfy the (anti-) commutation relations
\[
[J_{m}^{\Lambda_i}, J_{n}^{\Lambda_i}] = \sum_{c=1}^{2d} i f_{abc} J_{m+n}^{\Lambda_i} + \delta_{m+n,0} \delta_{a,b} m \ell 1_{\hat{\mathcal{H}}_{\lambda_i}},
\]
\[
[F_{m}^{\Lambda_i}, F_{n}^{\Lambda_i}] = \delta_{m+n,0} \delta_{a,b} 1_{\hat{\mathcal{H}}_{\lambda_i}},
\]
\[
[J_{m}^{\Lambda_i}, F_{n}^{\Lambda_i}] = 0.
\]
Here \( f_{abc} \) are the structure constants of \( g_C \oplus g_C \) with respect to a fixed orthonormal basis \( (e_a)_{a=1,\ldots,d} \) of the Lie algebra \( g_C \oplus 0 \) and \( (e_a)_{a=d+1,\ldots,2d} \) of \( 0 \oplus g_C \) satisfying the orthonormality condition with respect to the normalized Killing form
\[
-\frac{1}{2h^\vee} \text{tr} (\text{Ad}(e_a) \text{Ad}(e_b)) = \delta_{a,b} \quad a, b = 1, \ldots, 2d,
\]
where \( h^\vee \) is the dual Coxeter number of \( g_C \). Notice that \( f_{abc} = 0 \) if one of \( a, b, c \) lies in \( \{1, \ldots, d\} \) and another one in \( \{d+1, \ldots, 2d\} \). Moreover, we define
\[
J_{n}^{a}^{\Lambda_i} := -\frac{1}{2} \sum_{m \in \mathbb{Z}} \sum_{b,c=1}^{2d} f_{abc} F_{m}^{b} \Lambda_{i} F_{n-m}^{c} \Lambda_{i}.
\]
which has the commutation relations

\[ [J_{m}^{a,\Lambda_{i}}, J_{n}^{\delta,\Lambda_{i}}] = \sum_{c=1}^{2d} i f_{abc} J_{m+n+c}^{c,\pi_{R}} + \delta_{m+n,0} \delta_{a,\delta} m h^{\gamma} \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}}, \]

\[ [J_{m}^{a,\Lambda_{i}}, F_{n}^{\delta,\Lambda_{i}}] = \sum_{c=1}^{2d} i f_{abc} F_{m+n+c}^{c,\Lambda_{i}}. \]

Then the even operators \( L_{n}^{\Lambda_{i}} \) and odd \( G_{n}^{\Lambda_{i}} \), with \( n \in \mathbb{Z} \), defined through the super-Sugawara construction \[56\] (cf. also \[15\] Sect. 6, \[53\] Sect. 5.9) and \[87\] Sect. III.13) as

\[ G_{n}^{\Lambda_{i}} := \frac{1}{\sqrt{\ell + h^{\gamma}}} \sum_{a=1}^{2d} \sum_{m \in \mathbb{Z}} \left( J_{m}^{a,\Lambda_{i}} + \frac{1}{3} J_{m}^{a,\Lambda_{i}} \right) F_{n-m}^{a,\Lambda_{i}}, \]

\[ L_{n}^{\Lambda_{i}} := \sum_{a=1}^{2d} \left( \frac{1}{2(\ell + h^{\gamma})} \sum_{m \in \mathbb{Z}} : J_{m}^{a,\Lambda_{i}} J_{n-m}^{a,\Lambda_{i}} : - \frac{1}{2} \sum_{m \in \mathbb{Z}} m : F_{m}^{a,\Lambda_{i}} F_{n-m}^{a,\Lambda_{i}} : \right) + \frac{d}{8} \delta_{n,0} \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}}, \]

(3.11)

(where \( : \) stands for the normally ordered product) satisfy the Ramond super-Virasoro algebra (anti-) commutation relations

\[ [L_{m}^{\Lambda_{i}}, L_{n}^{\Lambda_{i}}] = (m - n)L_{m+n}^{\Lambda_{i}} + \frac{c}{12} (m^3 - m) \delta_{m+n,0} \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}}, \]

\[ [L_{m}^{\Lambda_{i}}, G_{n}^{\Lambda_{i}}] = \left( \frac{m}{2} - n \right) G_{m+n}^{\Lambda_{i}}, \]

\[ [G_{m}^{\Lambda_{i}}, G_{n}^{\Lambda_{i}}]_{+} = 2L_{m+n}^{\Lambda_{i}} + \frac{c}{3} (m^2 - \frac{1}{4}) \delta_{m+n,0} \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}}, \]

(3.12)

with central charge \( c = 2(\frac{d}{4} + \frac{d}{e + h^{\gamma}}) \[56\] (5.8)). Moreover, \( L_{0}^{\Lambda_{i}} = \hat{L}_{0}^{\Lambda_{i}} \) so that choosing \( D_{\lambda_{i}} := G_{0}^{\Lambda_{i}} \) proves the main part of our proposition.

Concerning the spectrum of \( D_{\lambda_{i}} \), we see from Eq. (3.9) that \( L_{0}^{\Lambda_{i}} \) is bounded below by \( h_{R} = \frac{d}{8} \). Hence, we find

\[ D_{2}^{\lambda_{i}} = L_{0}^{\Lambda_{i}} - \frac{c}{24} \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}} \geq \frac{d}{8} \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}} - \frac{1}{24} (d + \frac{2d}{\ell + h^{\gamma}}) \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}} = \frac{d}{12} - \left( \frac{\ell}{12(\ell + h^{\gamma})} \right) \mathbf{1}_{\hat{\mathcal{H}}_{\Lambda_{i}}} > 0, \]

as \( \ell \) and \( h^{\gamma} \) are positive. \( \square \)

As in the preceding proof, we denote the \emph{infinitesimal generator of rotations} in the representation \( \hat{\pi}_{\lambda_{i}} \) by \( L_{0}^{\Lambda_{i}} = \hat{L}_{0}^{\Lambda_{i}} \); it coincides then with \( D_{\lambda_{i}}^{2} \), up to an additive constant. It should be clear from the proof that the reason for constructing \( \hat{\pi}_{\lambda_{i}} \) as in (3.8) is that for the super-Sugawara construction, we need the number of fermionic fields to coincide with the dimension of the underlying Lie algebra, and in order for \( \hat{\pi}_{\lambda_{i}} \) to be graded, we need this dimension to be even. That is why we included an auxiliary copy of \( \mathfrak{g}_{\mathbb{C}} \). The representations of \( \hat{\mathfrak{g}}_{\mathbb{C}} \) we are actually interested are \( \lambda_{i} \), while \( \Lambda_{i} \) are the corresponding ones of \( \hat{\mathfrak{g}}_{\mathbb{C}} \oplus \hat{\mathfrak{h}}_{\mathbb{C}} \oplus \mathrm{CAR}(\mathcal{K}, \gamma) \) needed for the super-Sugawara construction.

**Proposition 3.2.** Let \( \hat{\pi}_{\mathrm{red}} := \bigoplus_{i=0}^{N-1} \hat{\pi}_{\lambda_{i}} \) and \( \delta_{\mathrm{red}} := \bigoplus_{i=0}^{N-1} \delta_{D_{\lambda_{i}}} = \delta_{D_{\mathrm{red}}} \), where \( D_{\mathrm{red}} := \bigoplus_{i=0}^{N-1} D_{\lambda_{i}} \). Define the compact differentiable subalgebra of \( B_{G_{t}} \) as

\[ \hat{\mathcal{R}}_{G_{t}}^{1} := \{ x \in \hat{\mathcal{R}}_{G_{t}} : \hat{\pi}_{\mathrm{red}}(x) \in \mathrm{dom}(\delta_{\mathrm{red}}), \delta_{\mathrm{red}}(\hat{\pi}_{\mathrm{red}}(x)) \in K(\hat{\mathcal{H}}_{\mathrm{red}}) \}. \]

The norm

\[ \| x \|_{1} = \| x \| + \| \delta_{\mathrm{red}}(\hat{\pi}_{\mathrm{red}}(x)) \|_{B(\hat{\mathcal{H}}_{\mathrm{red}})}, \quad x \in \hat{\mathcal{R}}_{G_{t}}^{1}, \]
is well-defined and turns $\mathcal{R}_{G_t}^1$ into a Banach algebra, so that the maps $\hat{\pi}_{\lambda_i} : \mathcal{R}_{G_t}^1 \to (\text{dom}(\delta_{\lambda_i}), \| \cdot \|_{B(\mathcal{R}_{\lambda_i})} + \|\delta_{\lambda_i}(\cdot)\|_{B(\mathcal{R}_{\lambda_i})})$, $i = 0, \ldots, N - 1$, are continuous.

**Proof.** \(\| \cdot \|_1\) is clearly a norm and well-defined on $\hat{\pi}^{-1}_{\text{red}}(\text{dom}(\delta_{\text{red}}))$ and turns it into a Banach algebra. Moreover, $(\mathcal{R}_{G_t}, \| \cdot \|)$ is a Banach algebra. As $\| \cdot \|_1$ is finer than $\| \cdot \|$, we see that the intersection $\mathcal{R}_{G_t} \cap \hat{\pi}^{-1}_{\text{red}}(\text{dom}(\delta_{\text{red}}))$ is a Banach algebra, too, w.r.t. $\| \cdot \|_1$. The fact that $(K(\mathcal{H}_{\text{red}}), \| \cdot \|_{B(\mathcal{H}_{\text{red}})})$ is complete shows then that $(\mathcal{R}_{G_t}^1, \| \cdot \|_1)$ is a Banach algebra.

**Proposition 3.3.** The Banach algebra $\mathcal{R}_{G_t}^1$ is stably-unital, and its finite-rank elements are dense.

**Proof.** We shall construct an approximate identity $(e_n)_{n \in \mathbb{N}}$ of finite-rank projections in $\mathcal{R}_{G_t}$. Once this is done, then given any $x \in \mathcal{R}_{G_t}^1$, we see that $e_n x$ is finite-rank and converges to $x$, proving also the second statement.

To start with, notice that $L_0^{\lambda_{\text{red}}} = \bigoplus_{p=0}^{N-1} L_0^{\lambda_p}$ has positive discrete spectrum with every eigenvalue of finite multiplicity $[72]$ (9.3.4]). The same holds for $L_0^{\pi_{R}}$ and hence for

$$L_0^{\lambda_{\text{red}}} := \bigoplus_{i=0}^{N-1} L_0^{\lambda_i} = L_0^{\lambda_{\text{red}}} \otimes 1_{\mathcal{H}_{\lambda_0}} \otimes 1_{\mathcal{H}_{\pi R}} + 1_{\mathcal{H}_{\lambda_{\text{red}}}} \otimes L_0^{\lambda_0} \otimes 1_{\mathcal{H}_{\pi R}} + 1_{\mathcal{H}_{\lambda_{\text{red}}}} \otimes 1_{\mathcal{H}_{\lambda_0}} \otimes L_0^{\pi_{R}}.$$

For every $n \in \mathbb{N}$, let $e_n$ denote the spectral projection of $L_0^{\lambda_{\text{red}}}$ corresponding to the eigenvalues less than $n$. It is a finite rank projection in $\mathcal{R}_{G_t}$. Similarly, let $\hat{e}_n \in B(\mathcal{H}_{\text{red}})$ denote the spectral projection of $L_0^{\lambda_{\text{red}}}$ corresponding to the eigenvalues less than $n + h_R$. Then $\hat{e}_n$ is a finite rank projections commuting with $D_{\text{red}}$ and hence $D_{\text{red}} \hat{e}_n$ is bounded with domain $\mathcal{H}_{\text{red}}$. In particular $\hat{e}_n \in \text{dom}(\delta_{\text{red}})$ and $\delta_{\text{red}}(\hat{e}_n) = 0$.

Moreover, $\hat{\pi}_{\text{red}}(e_n) = (1_{\mathcal{H}_{\lambda_{\text{red}}} \otimes q_{\mathcal{H}_{\lambda_0}} \otimes e_R}) \hat{e}_n$ is a subprojection of $\hat{e}_n$, and for every $x \in \mathcal{R}_{G_t}^1$, we have $\hat{\pi}_{\text{red}}(e_n x) = \hat{e}_n \hat{\pi}_{\text{red}}(x)$. It follows that $\pi_{\text{red}}(e_n)$ is compact and in $\text{dom}(\delta_{\text{red}})$, with $\delta_{\text{red}}(\pi_{\text{red}}(e_n))$ compact again so that $e_n \in \mathcal{R}_{G_t}^1$. In general, we furthermore have $\hat{\pi}_{\text{red}}(\mathcal{R}_{G_t}^1) \subset K(\mathcal{H}_{\text{red}})$. Recalling that $\delta_{\text{red}}(\hat{e}_n) = 0$ we have, for every $x \in \mathcal{R}_{G_t}^1$,

$$\|e_n x - x\|_1 = \|e_n x - x\| + \|\delta_{\text{red}}(\hat{e}_n \hat{\pi}_{\text{red}}(x)) - \delta_{\text{red}}(\hat{\pi}_{\text{red}}(x))\|_{B(\mathcal{H}_{\text{red}})}$$

$$= \|e_n x - x\| + \|\hat{e}_n \delta_{\text{red}}(\hat{\pi}_{\text{red}}(x)) - \delta_{\text{red}}(\hat{\pi}_{\text{red}}(x))\|_{B(\mathcal{H}_{\text{red}})}.$$

Now, $e_n \to 1_{\mathcal{H}_{\lambda_{\text{red}}}}$ and $\hat{e}_n \to 1_{\mathcal{H}_{\text{red}}}$ as $n \to \infty$, in the strong topology. As a consequence, $(e_n)_{n \in \mathbb{N}}$ is an approximate identity for $\mathcal{R}_{G_t}$ and $(\hat{e}_n)_{n \in \mathbb{N}}$ for $K(\mathcal{H}_{\text{red}})$. As $\delta_{\text{red}}(\pi_{\text{red}}(x)) = K(\mathcal{H}_{\text{red}})$ by assumption, the right hand side goes to zero. Thus $(e_n)_{n \in \mathbb{N}}$ is an approximate identity for $\mathcal{R}_{G_t}^1$. 

**Definition 3.4.** For every $i = 0, \ldots, N - 1$, the lowest energy projection $p_{\lambda_i}$ of $\lambda_i$ is defined to be the unique minimal projection in $B_{G_t}$ such that $\pi_{\lambda_i}(p_{\lambda_i}) = q_{\mathcal{H}_{\lambda_i}}$.

We recall that a subalgebra $A$ of a Banach algebra $B$ is said to be closed under holomorphic functional calculus if for every $x \in A$ and every function $f$ which is defined and holomorphic on a neighborhood of the spectrum $\sigma(x)$ of $x$ in the unitization $\hat{B}$, the element $f(x) \in \hat{B}$ lies in $\hat{A}$, cf. $[20]$ 3.App.C. For this it is sufficient to show that $(x - \mu 1)^{-1} \in \hat{A}$, for every $\mu \notin \sigma(x)$, owing to the Cauchy integral formula.
Proposition 3.5. The following holds:

(i) \( p_{\lambda_i} \in \hat{\mathfrak{R}}_{G_t} \), for every \( i = 0, \ldots, N - 1 \), determining a class \( [p_{\lambda_i}] \in K_0(\hat{\mathfrak{R}}_{G_t}) \).

(ii) \( \hat{\mathfrak{R}}_{G_t}^1 \subset \hat{\mathfrak{R}}_{G_t} \) is closed under holomorphic functional calculus and \( \overline{\hat{\mathfrak{R}}_{G_t}^1} = \hat{\mathfrak{R}}_{G_t} \).

(iii) \( K_0(\hat{\mathfrak{R}}_{G_t}^1) = K_0(\hat{\mathfrak{R}}_{G_t}) = \mathbb{Z}^N \).

Proof. (i) We see from the proof of Proposition 3.3 that the spectral projections \( e_n \) defined there are such that \( D_{\text{red}} \) is bounded with domain \( \hat{\mathfrak{H}}_{\text{red}} \) for every positive integer \( n \). As a consequence, for every \( x \in B_{G_t} \), \( \hat{\mathfrak{H}}_{\text{red}}(e_n x e_n) \) is in the domain of \( \delta_{\text{red}} \) and \( \delta_{\text{red}}(\hat{\pi}_{\text{red}}(e_n x e_n)) \) has finite rank so that \( e_n x e_n \in \hat{\mathfrak{R}}_{G_t}^1 \). For \( n \) sufficiently large we have that \( p_{\lambda_i} = e_n p_{\lambda_i} e_n \), thus \( p_{\lambda_i} \in \hat{\mathfrak{R}}_{G_t}^1 \).

(ii) Given \( x \in \hat{\mathfrak{R}}_{G_t}^1 \) and \( \mu \notin \sigma(x) \), where the spectrum is w.r.t. the C*-algebra \( \hat{\mathfrak{R}}_{G_t} \), we have to show that \( (x - \mu 1)^{-1} \in \hat{\mathfrak{R}}_{G_t}^1 \). To this end, notice that \( (x - \mu 1)^{-1} \in \hat{\mathfrak{R}}_{G_t} \cap \hat{\mathfrak{H}}_{\text{red}}(\text{dom}(\delta_{\text{red}})) \) because both \( \hat{\mathfrak{R}}_{G_t} \) and \( \hat{\mathfrak{H}}_{\text{red}}(\text{dom}(\delta_{\text{red}})) \) are closed under holomorphic functional calculus, cf. [7, Prop.3.2.29] for the second case, while for \( \hat{\mathfrak{R}}_{G_t} \) this is clear since it is a Banach algebra w.r.t. \( \| \cdot \| \). Moreover, the latter reference shows

\[
\delta_{\text{red}}(\hat{\pi}_{\text{red}}((x - \mu 1)^{-1})) = \delta_{\text{red}}(\hat{\pi}_{\text{red}}(x) - \mu 1)^{-1}
\]

which is compact because \( \delta_{\text{red}}(\hat{\pi}_{\text{red}}(x)) \) is compact. Hence, \( (x - \mu 1)^{-1} \in \hat{\mathfrak{R}}_{G_t}^1 \).

To see that \( \hat{\mathfrak{R}}_{G_t}^1 \) is a norm dense subalgebra of \( \hat{\mathfrak{R}}_{G_t} \) it is enough to note that

\[
\lim_{n \to \infty} \| e_n x e_n - x \| = 0
\]

for every \( x \in \hat{\mathfrak{R}}_{G_t} \).

(iii) follows immediately from (ii) together with [20, 3.App.C] or [3, 5.1.2].

Theorem 3.6. For every \( i = 0, \ldots, N - 1 \), \( (\hat{\mathfrak{R}}_{G_t}^1, \hat{\pi}_{\lambda_i}, D_{\lambda_i}) \) forms an even \( \theta \)-summable spectral triple with nonunital \( \hat{\mathfrak{R}}_{G_t}^1 \). It gives rise to the even JLO entire cyclic cocycle \( \tau_{\lambda_i} \in HE^e(\hat{\mathfrak{R}}_{G_t}^1) \). The JLO cocycle \( \tau_{\lambda_i} \) pairs with \( K_0(\hat{\mathfrak{R}}_{G_t}) \) and

\[
\tau_{\lambda_i}(p_{\lambda_j}) = \langle [\tau_{\lambda_i}], [p_{\lambda_j}] \rangle = \delta_{ij}, \quad i, j = 0, \ldots, N - 1.
\]

Proof. It is clear that \( \hat{\pi}_{\lambda_i} \) is a representation of \( \hat{\mathfrak{R}}_{G_t}^1 \) with image in \( \text{dom}(\delta_{D_{\lambda_i}}) \) – in fact, in the even part of \( \text{dom}(\delta_{D_{\lambda_i}}) \) because \( 1_{\mathfrak{H}_{\lambda_i}} \otimes \mathfrak{H}_{\sigma_R} \otimes e_R \) is even and \( B(\mathfrak{H}_{\lambda_i}) \otimes 1_{\mathfrak{H}_{\sigma}} \otimes 1_{\mathfrak{H}_{\sigma_R}} \) is even. As explained in [15, Sect.6] (cf. also [10, III.7.(iv)] using the character formulae in [53, Sect.10]), both \( e^{-tL_{\alpha_i}} \) and \( e^{-t^2L_{\alpha_i}} \) on \( \mathfrak{H}_{\lambda_i} \) and \( \mathfrak{H}_{\sigma_R} \), respectively, are trace-class, thus \( e^{-tL_{\alpha_i}} \) on \( \mathfrak{H}_{\lambda_i} \), for every \( t > 0 \), and we have an even \( \theta \)-summable spectral triple. It induces an even JLO cocycle, which pairs with \( K_0(\hat{\mathfrak{R}}_{G_t}) \) according to Theorem 2.4(ii) and Proposition 3.5(ii)&(iii). The actual values can be computed as follows: according to Definition 3.3 \( \hat{\pi}_{\lambda_i}(p_{\lambda_j}) = 0 \) if \( i \neq j \), and

\[
p := \hat{\pi}_{\lambda_i}(p_{\lambda_i}) = q_{\lambda_i} \otimes q_{\lambda_j} \otimes e_R,
\]

which yields \( pD_{\lambda_i}p = 0 \), \( \dim(p\mathfrak{H}_{\lambda_i}^+ \cap \mathfrak{H}_{\lambda_i}^-) = 1 \) and \( \dim(p\mathfrak{H}_{\lambda_i}^+ \cap \mathfrak{H}_{\lambda_i}^-) = 0 \), so

\[
\tau_{\lambda_i}(p_{\lambda_j}) = \langle [\tau_{\lambda_i}], [p_{\lambda_j}] \rangle = \text{ind}_{\hat{\pi}_{\lambda_i}^+} (pD_{\lambda_i}p) = 1 - 0 = 1.
\]

We would like to look at this pairing from the point of view of Kasparov’s KK-theory, for which we are now providing a few preliminaries based on the summary in [11]. A thorough introduction with proofs can be found in [3, Ch.17&18].
Given two stably-unital separable C*-algebras \( A, B \), a Kasparov \((A, B)\)-module is a tuple \((E, \phi, F)\), where \( E \) is a countably generated \( \mathbb{Z}_2 \)-graded Hilbert \( B \)-module, \( \phi : A \to B(E) \) is a graded *-homomorphism, and \( F \in B(E) \) has degree one, such that
\[
(F - F^*) \phi(a), \quad (F^2 - 1) \phi(a), \quad [F, \phi(a)]
\] (3.15)
lie all in the compact operators \( K(E) \) (the norm closure of finite-rank operators) on \( E \), for all \( a \in A \). With a suitable concept of homotopy, one defines \( KK(A, B) \) as the set of homotopy equivalence classes of Kasparov \((A, B)\)-modules. E.g. for any unitary \( u \in B(E) \) of degree zero, \((E, Ad(u) \circ \phi, Ad(u)(F))\) is again a Kasparov \((A, B)\)-module, which is homotopy equivalent to \((E, \phi, F)\), and therefore the two give rise to the same element in \( KK(A, B) \). We write \( \hat{H}_A \) for the standard right Hilbert \( A \)-module \( A \otimes l^2(\mathbb{Z}) \) and \( \hat{H}_A \) for the corresponding \( \mathbb{Z}/2 \)-graded one \( \hat{H}_A \oplus \hat{H}_A \) with grading \( 1 \oplus -1 \).

For our immediate purposes here and in the following section, the most relevant facts about KK-theory can be summarized as follows, with \( A, B, C \) stably-unital separable C*-algebras:

1. There is a direct sum for Kasparov \((A, B)\)-modules, which passes to the quotient \( KK(A, B) \) and turns \( KK(A, B) \) into an abelian group.
2. There is a canonical identification of \( KK(C, A) \) with the K-theory group \( K_0(A) \) (as additive groups). Similarly, there is a canonical identification of \( KK(A, C) \) with the K-homology group \( K^0(A) \).

This identification works as follows for \( K_0(A) \):
\[
[p_+] - [p_-] \in K_0(A) \mapsto \left[ \hat{H}_A, \phi_{p_+} \oplus \phi_{p_-}, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \in KK(C, A),
\]

as any element in \( K_0(A) \) may be written as a formal difference \([p_+] - [p_-]\), with (not unique) \( p_+, p_- \in A \otimes K \), and \( \phi_p : \mathbb{C} \to \mathbb{C} \mapsto tp \in A \otimes K \), for any projection \( p \in A \otimes K \). Owing to this identification, we may consider \([p]\) as an element of \( KK(C, A) \), for every \( p \in A \otimes K \), cf. [3, 17.5-17.6]

On the other hand, \( K^0(A) \) is by definition the group generated by homotopy classes of even Fredholm modules on a standard \( \mathbb{Z}_2 \)-graded separable Hilbert space \( \hat{H} \), i.e., classes \((\hat{H}, \phi, F)\) with \( \phi : A \to B(\hat{H}) \) a graded *-homomorphism and \( F \) a unitary selfadjoint operator on \( B(\hat{H}) \) of degree 1 such that the graded commutator satisfies \([F, \phi(a)]_{(+)} \in K(\hat{H})\), for all \( a \in A \), cf. [20, p.294]. Thus \( K^0(A) \) coincides with \( KK(A, C) \).

3. Every *-homomorphism \( \phi : A \to B \) naturally defines a \( KK(A, B) \)-element \([\phi]\) as the homotopy class of \((B, \phi, 0)\), where we have identified \( B(\hat{H}) \) with the multiplier algebra \( M(B) \) of \( B \). \([\phi]\) depends only on the unitary-equivalence class in \( M(B) \), cf. [3, 17.1-17.2].

4. There exists a well-defined bilinear map \( \times \), the so-called Kasparov product
\[
KK(A, B) \times KK(B, C) \to KK(A, C),
\]

which is associative. It is in general complicated to define and we refer to [3, Ch.18], but for the following special cases we provide formulae.

5. Suppose \( A, B, C \) are trivially graded and given two classes of Kasparov modules
\[
\left[ \hat{H}_B, \psi \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \in KK(A, B), \quad \left[ \hat{H}_C, \phi \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \in KK(B, C),
\]

with \( \psi \) a *-homomorphism from \( A \) to \( B \otimes K \) and \( \phi \) is a *-homomorphism from \( B \) to \( C \otimes K \); the latter extends to \( B \otimes K \), denoted again by \( \phi \). As explained in [3, Ex.18.4.2(c)], the Kasparov product of the two is then given by the element
\[
\left[ \hat{H}_B, \psi \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \times \left[ \hat{H}_C, \phi \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] = \left[ \hat{H}_C, \phi \circ \psi \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right]
\]
in \( KK(A,C) \).

(6) If \( \psi : A \to B \) and \( \phi : B \to C \) are \(*\)-homomorphisms then

\[
\{\psi\} \times \{\phi\} = \{\phi \circ \psi\}. \tag{3.16}
\]

If \( id_A \) is the identity automorphism of \( A \) then \( \{id_A\} \) is the neutral element in \( KK(A,A) \) for the Kasparov product. Hence \( KK(A,A) \) is a unital ring, cf. [3 18.7.1].

The Kasparov product \( KK(\mathbb{C},A) \times KK(\mathbb{C},A) \to KK(\mathbb{C},\mathbb{C}) \simeq \mathbb{Z} \) gives rise to an index pairing between \( K_0(A) = KK(\mathbb{C},A) \) and \( K^0(\mathbb{A}) = KK(\mathbb{A},\mathbb{C}) \) and to a corresponding map \( \gamma_\alpha : K^0(\mathbb{A}) \to \text{Hom}(K_0(\mathbb{A}),\mathbb{Z}) \). Moreover, the Kasparov product \( KK(\mathbb{C},A) \times KK(\mathbb{C},A) \to KK(\mathbb{C},\mathbb{C}) = K_0(\mathbb{A}) \) gives rise to a map \( \gamma_\beta : KK(A,A) \to \text{End}(K_0(\mathbb{A})) \). Similarly, the Kasparov product \( KK(A,A) \times KK(A,C) \to KK(A,C) = K^0(A) \) gives rise to a map \( \gamma_\epsilon : KK(A,A) \to \text{End}(K^0(\mathbb{A})) \).

The C*-algebra \( \mathfrak{R}_{G_i} \) lies in the so called bootstrap class [3 V.1.5.4]. Hence we can apply the universal coefficient theorem [3 V.1.5.8] which implies that \( \gamma_\beta : KK(\mathfrak{R}_{G_i},\mathfrak{R}_{G_i}) \to \text{End}(K_0(\mathfrak{R}_{G_i})) \) is an isomorphism (because \( \text{Ext}_{\mathfrak{R}_{G_i}}^1(\mathbb{Z},\mathbb{Z}) = 0 \)) and hence, using Proposition [3.1], \( KK(\mathfrak{R}_{G_i},\mathfrak{R}_{G_i}) \simeq \text{End}(K_0(\mathfrak{R}_{G_i})) \simeq \text{End}(\mathbb{Z}) \). It also implies that \( \gamma_\alpha : K^0(\mathfrak{R}_{G_i}) \to \text{Hom}(K_0(\mathfrak{R}_{G_i}),\mathbb{Z}) \) is an isomorphism so that \( K^0(\mathfrak{R}_{G_i}) \simeq \mathbb{Z} \). It follows that \( \gamma_\epsilon : KK(\mathfrak{R}_{G_i},\mathfrak{R}_{G_i}) \to \text{End}(K^0(\mathfrak{R}_{G_i})) \) is an isomorphism too so that \( KK(\mathfrak{R}_{G_i},\mathfrak{R}_{G_i}) \simeq \text{End}(K^0(\mathfrak{R}_{G_i})) \simeq \text{End}(\mathbb{Z}) \).

For any \( i = 0, \ldots, N-1 \), one associates to the \( \delta \)-summable spectral triple \( (\mathfrak{R}_{G_i}^0, \hat{\tau}_\lambda^i, D_{\lambda_i}) \) determining \( \tau_\lambda^i \) the Fredholm module \( (\hat{\mathcal{H}}_{\lambda_i}, \hat{\pi}_\lambda, \text{sgn} D_{\lambda_i}) \). Here \( \text{sgn} D_{\lambda_i} = D_{\lambda_i}/|D_{\lambda_i}|^{-1} \) is the signature of \( D_{\lambda_i} \) (recall that 0 is not in the spectrum of \( D_{\lambda_i} \) according to Proposition [3.1]). It gives rise to the same index map as \( D_{\lambda_i} \), cf. [20 Sect. IV.8.δ]. We write \( \varepsilon_{\lambda_i} := [\hat{\mathcal{H}}_{\lambda_i}, \hat{\pi}_\lambda, \text{sgn} D_{\lambda_i}] \in K^0(\mathfrak{R}_{G_i}) \) for the corresponding K-homology class. The following proposition is included for the sake of completeness and to make things as explicit as possible for our setting.

**Proposition 3.7.** The classes \( \varepsilon_{\lambda_i} \) and \( [p_{\lambda_i}] \), with \( i = 0, \ldots, N-1 \), generate \( K^0(\mathfrak{R}_{G_i}) \) and \( K_0(\mathfrak{R}_{G_i}) \), respectively. For every \( i, j = 1, \ldots, N \), we have

\[
[p_{\lambda_i}] \times \varepsilon_{\lambda_i} = \tau_{\lambda_i}(p_{\lambda_j}) = \delta_{i,j}. \tag{3.17}
\]

In other words, the pairing of the JLO cocycle with \( K_0(\mathfrak{R}_{G_i}) \) is given by the Kasparov product between the corresponding K-homology class and \( K_0(\mathfrak{R}_{G_i}) \).

**Proof.** The fact that \( F := \text{sgn} D_{\lambda_i} \) has degree 1 with \( F^2 = 1 \) and \( F^* = F = F^{-1} \) and that \( \hat{\pi}_\lambda \) has support in \( \hat{H}_{\lambda_i,\epsilon} \) implies that we may write \( F = \begin{pmatrix} 0 & W^* \\ W & 0 \end{pmatrix} \) with a certain unitary \( W \). Then

\[
\varepsilon_{\lambda_i} = \left[ \hat{\mathcal{H}}_{\lambda_i}, \begin{pmatrix} \hat{\pi}_{\lambda_i,\epsilon} & 0 \\ 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & W^* \\ W & 0 \end{pmatrix} \right] = \left[ \hat{\mathcal{H}}_{\lambda_i}, \begin{pmatrix} W\hat{\pi}_{\lambda_i,\epsilon}(\cdot)W^* & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \right],
\]

where the last equality is a consequence of unitary equivalence via the unitary \( \begin{pmatrix} W & 0 \\ 0 & 1 \end{pmatrix} \in B(\hat{\mathcal{H}}_{\lambda_i}) \).
Now the Kasparov product of this element with \([p_{\lambda_j}]\) can be calculated using property (5) above, as the involved algebras \(\mathbb{C}\) and \(\mathcal{R}_{G_t}\) are trivially graded. Thus

\[
\begin{bmatrix}
\mathcal{H}_{\mathcal{R}_{G_t}}\otimes \phi_{p_{\lambda_j}} + 0, \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \\
\end{bmatrix} \times \begin{bmatrix}
\mathcal{H}_{\mathcal{H}_0}, W\tilde{\pi}_{\lambda_i,+}(\cdot)W^* \oplus 0, \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \\
\end{bmatrix} = \begin{bmatrix}
\mathcal{H}_{\mathcal{H}_0}, W\tilde{\pi}_{\lambda_i,+}(\phi_{p_{\lambda_j}}(\cdot))W^* \oplus 0, \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \\
\end{bmatrix} = \delta_{i,j} \begin{bmatrix}
\mathcal{H}, \phi_{WpW^*} + 0, \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \\
\end{bmatrix}
\]

as product \(KK(\mathbb{C}, \mathcal{R}_{G_t}) \times KK(\mathcal{R}_{G_t}, \mathbb{C}) \to KK(\mathbb{C}, \mathbb{C})\). Here we used the identification of \(\mathcal{H}_{\mathcal{H}_i}\) with the standard \(\mathbb{Z}_2\)-graded Hilbert space and Hilbert \(\mathbb{C}\)-module \(\mathcal{H}\) and the definition of \(p_{\lambda_i}\) and \(\tilde{\pi}_{\lambda_i,+}\) in the last line, where \(p\) as in (3.14). Now \(p\) and hence \(WpW^*\) are rank one projections and therefore \(\mathcal{H}, \phi_{WpW^*} + 0, \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \) is the generator of \(KK(\mathbb{C}, \mathbb{C}) \simeq \mathbb{Z}\). Thus

\[
\begin{bmatrix}
\mathcal{H}_{\mathcal{R}_{G_t}}\otimes \phi_{p_{\lambda_j}} + 0, \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \\
\end{bmatrix} \times \begin{bmatrix}
\mathcal{H}_{\mathcal{H}_0}, W\tilde{\pi}_{\lambda_i,+}(\cdot)W^* \oplus 0, \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \\
\end{bmatrix} = \delta_{i,j}, \quad i, j = 0, \ldots, N - 1.
\]

As \(K_0(\mathcal{R}_{G_t}) \simeq \mathbb{Z}^n\) and \(K^0(\mathcal{R}_{G_t}) \simeq \mathbb{Z}^n\), we therefore see that the elements \([p_{\lambda_i}]\) and \([\varepsilon_{\lambda_i}]\), with \(i, j = 0, \ldots, N - 1\), generate \(K_0(\mathcal{R}_{G_t})\) and \(K^0(\mathcal{R}_{G_t})\) respectively. Together with Theorem 3.6 we get

\[ [p_{\lambda_i}] \times [\varepsilon_{\lambda_i}] = \delta_{i,j} = \tau_{\lambda_i}(p_{\lambda_j}). \]

\[ \square \]

Since \([p_{\lambda_i}]\) and \([\varepsilon_{\lambda_i}]\), with \(i = 0, \ldots, N - 1\), generate \(K_0(\mathcal{R}_{G_t})\) and \(K^0(\mathcal{R}_{G_t})\), respectively, the maps \([\lambda_i] \mapsto [p_{\lambda_i}]\) and \([\lambda_i] \mapsto [\varepsilon_{\lambda_i}]\), \(i = 0, \ldots, N - 1\), give rise to (surjective) group isomorphisms

\[
\phi_{-1} : R^\ell(\mathcal{L}G) \to K_0(\mathcal{R}_{G_t})
\]

and

\[
\phi_1 : R^\ell(\mathcal{L}G) \to K^0(\mathcal{R}_{G_t}).
\]

Now the regular representation of \(R^\ell(\mathcal{L}G)\) together with the map \(\phi_1\) induces an injective ring homomorphism of \(R^\ell(\mathcal{L}G)\) into \(\text{End}(K^0(\mathcal{R}_{G_t}))\) and hence an injective ring homomorphism

\[
\phi_0 : R^\ell(\mathcal{L}G) \to KK(\mathcal{R}_{G_t}, \mathcal{R}_{G_t})
\]

determined by the condition

\[
\phi_0([\lambda_i]) \times [\varepsilon_{\lambda_j}] = \sum_{k=0}^{N-1} \mathcal{N}^k_{ij} [\varepsilon_{\lambda_k}] \quad i, j \in \{0, \ldots, N - 1\} \quad (3.18)
\]

Then, using Eqs. (3.5), (3.6), (3.7) and (3.17), it is straightforward to see that

\[
[p_{\lambda_i}] \times \phi_0([\lambda_i]) = \sum_{k=0}^{N-1} \mathcal{N}^k_{ij} [p_{\lambda_k}] \quad i, j \in \{0, \ldots, N - 1\}. \quad (3.19)
\]

For

\[
x = \sum_{i=0}^{N-1} m_i [\lambda_i] \in R^\ell(\mathcal{L}G)
\]
we define \( \bar{x} \in R^e(LG) \) by
\[
\bar{x} = \sum_{i=0}^{N-1} m_i[\lambda_i].
\]
Using this notation we see that for \( x, y \in R^e(LG) \) we have
\[
\phi_0(x) \times \phi_1(y) = \phi_1(xy), \quad \phi_{-1}(y) \times \phi_0(x) = \phi_{-1}(\bar{x}y).
\]

We resume the above discussion in the following theorem.

**Theorem 3.8.** There exist necessarily unique group isomorphisms
\[
\phi_{-1} : R^e(LG) \to K_0(\mathfrak{R}_{G_{\ell}})
\]
\[
\phi_1 : R^e(LG) \to K^0(\mathfrak{R}_{G_{\ell}})
\]
such that \( \phi_{-1}([\lambda_i]) = [p_{\lambda_i}] \) and \( \phi_1([\lambda_i]) = \varepsilon_{\lambda_i}, i = 0, \ldots, N - 1. \) Moreover there exists a unique ring homomorphism
\[
\phi_0 : R^e(LG) \to KK(\mathfrak{R}_{G_{\ell}}, \mathfrak{R}_{G_{\ell}})
\]
such that \( \phi_0(x) \times \phi_1(y) = \phi_1(xy) \) for all \( x, y \in R^e(LG). \) \( \phi_0 \) is injective and satisfies
\[
\phi_{-1}(y) \times \phi_0(x) = \phi_{-1}(\bar{x}y)
\]
for all \( x, y \in R^e(LG). \) Moreover,
\[
\phi_{-1}([\lambda_i]) \times \phi_0([\lambda_i]) \times \phi_1([\lambda_j]) = N^k_{ij}
\]
for all \( i, j, k \in \{0, \ldots, N - 1\}. \)

**Remark 3.9.** The maps \( \phi_1 \) and \( \phi_{-1} \) can be recovered from \( \phi_0 \) through the identities
\[
\phi_1(x) = \phi_0(x) \times \varepsilon_{\lambda_0}, \quad \phi_{-1}(x) = [p_{\lambda_0}] \times \phi_0(\bar{x}), \quad x \in R^e(LG).
\]

4. **DHR FUSION RING AND CONFORMAL NETS**

Conformal nets describe chiral conformal CFTs in the operator algebraic approach to quantum field theory [49]. They can be considered as a functional analytic analogue of vertex algebras [39, 55], see [17].

Let \( \mathcal{I} \) denote the set of nondense nonempty open intervals in \( S^1 \) and \( Diff(S^1) \) the infinite-dimensional Lie group of orientation-preserving smooth diffeomorphisms of \( S^1 \) [71]. \( Diff(S^1) \) contains the group of Möbius transformations of \( S^1 \) which is isomorphic to \( PSL(2, \mathbb{R}) \). Accordingly we will consider \( PSL(2, \mathbb{R}) \) as a (Lie) subgroup of \( Diff(S^1) \).

For \( I \in \mathcal{I} \) we denote by \( I^c \) the interior of the complement of \( I \), which lies again in \( \mathcal{I} \). A local conformal net \( \mathcal{A} \) over \( S^1 \) (cf. [40, 48]) consists of a family of von Neumann algebras \( (\mathcal{A}(I))_{I \in \mathcal{I}} \) acting on a common separable Hilbert space \( \mathcal{H} \) together with a given strongly continuous unitary representation \( U \) of \( PSL(2, \mathbb{R}) \) on \( \mathcal{H} \) satisfying

- **isotopy:** \( \mathcal{A}(I_1) \subset \mathcal{A}(I_2) \) if \( I_1 \subset I_2 \), for \( I_1, I_2 \in \mathcal{I} \);
- **locality:** elements of \( \mathcal{A}(I_1) \) commute with those of \( \mathcal{A}(I_2) \) whenever \( I_1 \cap I_2 = \emptyset \);
- **covariance:** \( U(\gamma)\mathcal{A}(I)U(\gamma)^* = \mathcal{A}(\gamma I) \) for all \( \gamma \in PSL(2, \mathbb{R}) \) and \( I \in \mathcal{I} \);
- **positivity of the energy:** the conformal Hamiltonian \( L_0 \), defined by the equation
  \[
  U(R_\alpha) = e^{i\alpha L_0} \quad (\forall \alpha \in \mathbb{R}),
  \]
  is positive, where \( (R_\alpha)_{\alpha \in \mathbb{R}} \subset PSL(2, \mathbb{R}) \) stands for the rotation subgroup;
- **existence, uniqueness and cyclicity of the vacuum:** up to phase there exists a unique unit vector \( \Omega \in \mathcal{H} \) called the “vacuum vector” which is invariant under the action of \( U \); moreover, it is cyclic for the von Neumann algebra \( \bigvee_{I \in \mathcal{I}} \mathcal{A}(I) \);
- **diffeomorphism covariance:** \( U \) extends (uniquely) to a strongly continuous projective unitary representation of \( Diff(S^1) \) denoted again by \( U \) and satisfying
  \[
  U(\gamma)\mathcal{A}(I)U(\gamma)^* = \mathcal{A}(\gamma I),
  \]
  \[
  \gamma |_{I = I} = \text{id}_I \Rightarrow \text{Ad}(U(\gamma)) |_{\mathcal{A}(I)} = \text{id}_{\mathcal{A}(I)},
  \]
  for all \( \gamma \in Diff(S^1) \) and \( I \in \mathcal{I} \).
There are many known important consequences of the above definition, cf. [48] and references therein for a collection with proofs. We shall need the following three: $\mathcal{A}(I)' = \mathcal{A}(I')$, for every $I \in \mathcal{I}$ (Haag duality), and $\bigvee_{I \in \mathcal{I}} \mathcal{A}(I) = B(\mathcal{H})$ (irreducibility), $\mathcal{A}(I)$ is a factor, for every $I \in \mathcal{I}$ (factoriality).

A conformal net is said to have the split property if for any inclusion $I_- \subset I_+$, there is a type I factor $F$ such that $\mathcal{A}(I_-) \subset F \subset \mathcal{A}(I_+)$. By [24, Thm. 3.2] if a conformal net $\mathcal{A}$ satisfies the trace class condition, i.e. if $\text{Tr}(q^{L_0}) < +\infty$ for all $q \in (0,1)$ then $\mathcal{A}$ also satisfies the split property.

A representation of $\mathcal{A}$ is a family $\pi = (\pi_I)_{I \in \mathcal{I}}$ of (unital) $*$-representations $\pi_I$ of $\mathcal{A}(I)$ on a common Hilbert space $\mathcal{H}_\pi$ such that $\pi_{I_2} |_{\mathcal{A}(I_1)} = \pi_{I_1}$ whenever $I_1 \subset I_2$. The representation $\pi$ is called locally normal if $\pi_I$ is normal for every $I \in \mathcal{I}$; this is always the case if $\mathcal{H}_\pi$ is separable.

A locally normal representation $\pi$ of the conformal net $\mathcal{A}$ is always (Möbius) covariant with positive energy in the sense that there exists a unique strongly continuous unitary representation $U_\pi$ of the universal covering $\text{PSL}(2,\mathbb{R})^{(\infty)}$ such that

$$U_\pi(g)\pi_I(a)U_\pi(g)^* = \pi_{gI}(U(g)aU(g)^*), \quad g \in \text{PSL}(2,\mathbb{R})^{(\infty)} \ a \in \mathcal{A}(I),$$

where $\hat{g}$ is the image of $g$ in $\text{PSL}(2,\mathbb{R})$ under the covering map, and such that $U_\pi(g) \in \bigvee_{I \in \mathcal{I}} \pi_I(\mathcal{A}(I))$, for all $g \in \text{PSL}(2,\mathbb{R})^{(\infty)}$, cf. [23]. Moreover, the infinitesimal generator $L_0^\pi$ of the lifting of the rotation subgroup turns out to be positive [89].

The vacuum representation $\pi_0$ on the vacuum Hilbert space $\mathcal{H}_{\pi_0} := \mathcal{H}$ is defined by $\pi_0|_1(x) = x$, for all $I \in \mathcal{I}$ and all $x \in \mathcal{A}(I)$, and it is obviously locally normal. The unitary equivalence class of a locally normal representation $\pi$ up to unitary equivalence is called a (DHR) sector and denoted by $[\pi]$. If $\pi$ is irreducible then $[\pi]$ is said to be an irreducible sector.

A covariant representation $\pi$ on the vacuum Hilbert space $\mathcal{H}$ is said to be localized in $I_0$ if $\pi_{I'} = \text{id}$, for all $I \supset I_0$. In this case, we have $\pi_I(\mathcal{A}(I)) \subset \mathcal{A}(I)$, for all $I \in \mathcal{I}$ containing $I_0$, i.e. $\pi_I$ is an endomorphism of $\mathcal{A}(I)$. If $\pi$ is any representation of $\mathcal{A}$ on a separable Hilbert space $\mathcal{H}_\pi$ and $I_0$ is any interval in $\mathcal{I}$ then there exists a covariant representation localized in $I_0$ and unitarily equivalent to $\pi$.

The universal $C^*$-algebra [33, Sect.5.3] of $\mathcal{A}$ is the unique (up to isomorphism) unital $C^*$-algebra $C^*(\mathcal{A})$ such that

- for every $I \in \mathcal{I}$, there are unital embeddings $\iota_I : \mathcal{A}(I) \to C^*(\mathcal{A})$, such that $\iota_{I_2} |_{\mathcal{A}(I_1)} = \iota_{I_1}$ if $I_1 \subset I_2$, and all $\iota_I(\mathcal{A}(I)) \subset C^*(\mathcal{A})$ together generate $C^*(\mathcal{A})$ as $C^*$-algebra;
- for every representation $\pi$ of $\mathcal{A}$ on $\mathcal{H}_\pi$, there is a unique representation $\hat{\pi} : C^*(\mathcal{A}) \to B(\mathcal{H}_\pi)$ such that $\pi_I = \hat{\pi} \circ \iota_I$, $I \in \mathcal{I}$,

cf. also [12] [32] [47]. If $\pi$ is locally normal then we say that $\hat{\pi}$ is locally normal. If $\pi$ is localized in $I_0$ then we say that $\hat{\pi}$ is localized in $I_0$. In the following, we shall denote $\hat{\pi}$ simply by $\pi$ since it will be clear from the context whether $\pi$ is a representation of $\mathcal{A}$ or the corresponding representation of $C^*(\mathcal{A})$.

There is a natural correspondence between localized representations and covariant localized endomorphisms of $C^*(\mathcal{A})$. Namely, for every locally normal representation $\pi$ of $C^*(\mathcal{A})$ localized in $I_0$, there is an endomorphism $\rho$ of $C^*(\mathcal{A})$ which is, in an appropriate sense, covariant and localized in $I_0$, such that $\pi = \pi_0 \circ \rho$ [33] [17], cf. also [12, Sect.2].

If $\pi$ is a representation of $\mathcal{A}$ on a separable Hilbert space then, for any $I \in \mathcal{I}$, $\pi_{I'}(\mathcal{A}(I')')$ and $\pi_I(\mathcal{A}(I))$ are factors and, as a consequence of locality, $\pi_I(\mathcal{A}(I)) \subset \pi_{I'}(\mathcal{A}(I')')$. The square of the index $[\pi_I(\mathcal{A}(I')) : \pi_I(\mathcal{A}(I))] \in [1, +\infty]$ of the subfactor
the Grothendieck ring) and call it the DHR fusion ring \( \mathcal{R}_A \). As in [12] we denote by \( \hat{\text{fusion}} \) semiring \( A \). Complete rationality appears to be the right notion for rational chiral CFTs in the \( \mathbb{R} \) case, with finite statistical dimension. If it satisfies the split property and it has only finitely many irreducible sectors, the set of equivalence classes of locally normal representations of \( \mathcal{A} \) with finite statistical dimension becomes a commutative unital semiring (without 0) \( \mathcal{R}_A \) called the DHR fusion semiring \([73]\). As in [12] we denote by \( \tilde{\mathcal{R}}_A \) the corresponding ring of formal differences (Grothendieck ring) and call it the DHR fusion ring of \( \mathcal{A} \).

An important class of nets is the class of completely rational nets introduced in [61]. Complete rationality appears to be the right notion for rational chiral CFTs in the operator algebraic setting of conformal nets. By \([70]\) a net \( \mathcal{A} \) is completely rational if and only if it satisfies the split property and it has only finitely many irreducible sectors, all with finite statistical dimension. If \( \mathcal{A} \) is completely rational then \( \tilde{\mathcal{R}}_A \) is finitely generated. We shall henceforth restrict ourselves to the case where \( \mathcal{A} \) is completely rational.

In [12] Sect.3 a locally normal universal C*-algebra is constructed for general \( \mathcal{A} \) which is more manageable than \( C^*(\mathcal{A}) \) and which can be used as a substitute for it if only locally normal representations are considered, as typically the case. As \( \mathcal{A} \) is assumed here to be completely rational with \( N < \infty \) irreducible sectors, we may however work with a simpler C*-algebra isomorphic to the locally normal C*-algebra and also defined in [12]. It is the reduced locally normal universal C*-algebra, defined as follows. For each sector of \( \mathcal{A} \) consider a fixed representative representation of \( C^*(\mathcal{A}) \) and let \( (\pi_{\text{red}}, \mathcal{H}_{\text{red}}) \) denote their direct sum and \( C^*_{\text{red}}(\mathcal{A}) := \pi_{\text{red}}(C^*(\mathcal{A})) \), which turns out to be isomorphic to \( B(\mathcal{H})^{\boxtimes N} \) [12, Thm.3.3]. Then every locally normal representation \( \pi \) of \( C^*(\mathcal{A}) \) is quasi-equivalent to a subrepresentation of \( \pi_{\text{red}} \) and hence gives rise to a unique normal representation \( \pi' \) of \( C^*_{\text{red}}(\mathcal{A}) \) such that \( \pi' \circ \pi_{\text{red}} = \pi \). As a consequence, for every covariant localized endomorphism \( \rho \) of \( C^*(\mathcal{A}) \) there exists a unique endomorphism \( \hat{\rho} \) of \( C^*_{\text{red}}(\mathcal{A}) \) such that \( \hat{\rho}(\pi_{\text{red}}(x)) = \pi_{\text{red}}(\rho(x)) \) for all \( x \in C^*(\mathcal{A}) \). If \( \pi_0 \circ \rho \) has finite statistical dimension then \( \hat{\rho} \) restricts to an endomorphism of the compact universal C*-algebra

\[
\mathcal{R}_A := C^*_{\text{red}}(\mathcal{A}) \cap K(\mathcal{H}_{\text{red}}),
\]

introduced in [12, Prop.3.4].

Let us now apply the general theory to our setting of loop groups. Let \( G \) be a connected simply connected compact simple Lie group, with the same notation as introduced in Section 5. The loop group net of \( G \) at level \( \ell \) on \( \mathcal{H} := \mathcal{H}_{\lambda_0} \) is defined as

\[
\mathcal{A}_{G_{\ell}}(I) := \{ \lambda_0(g) : g \in L_I G \}, \quad I \in \mathcal{I},
\]

where the local subgroups \( L_I G \subset L G, I \in \mathcal{I} \), are defined by

\[
L_I G := \{ g \in L G : g |_{I'} = 1 \},
\]

cf. [40]. It is a conformal net satisfying the split property. The locality property of the net can be proved in various ways. In particular, it follows from the following lemma that we will also use later, cf. [77, Prop.1.1.2]. The fact that \( G \) is simply connected is crucial here.

**Lemma 4.1.** Let \( \lambda : G \to U(\mathcal{H}_{\lambda})/U(1) \) be a strongly continuous projective unitary representation of \( L G \), let \( g \mapsto \hat{\lambda}(g) \in U(\mathcal{H}_{\lambda}) \) be any given choice of the representatives of \( \lambda(g) \) and let \( I \in \mathcal{I} \). Then every unitary operator \( u_I \) satisfying \( u_I \lambda(g) u_I^* = \lambda(g) \), for all \( g \in L_I G \), commutes with \( \hat{\lambda}(g) \), for all \( g \in L_I G \).
The projective equality $u_I\lambda(g)u_I^* = \lambda(g)$, $g \in L_I G$ implies that the map $L_I G \ni g \mapsto \chi(g) \in \mathbb{U}(1)$ defined by
\begin{equation}
\chi(g)1_{\mathcal{H}_\lambda} := u_I\hat{\lambda}(g)u_I^*\hat{\lambda}(g)^*, \quad g \in L_I G,
\end{equation}
is a continuous character of $L_I G$. But $L_I G$ is a perfect group by [77] Lemma 1.1.1 and hence $\chi(g) = 1$ for all $g \in L_I G$.

Let $\lambda_i, i = 0, \ldots, N-1$, be the irreducible level $\ell$ positive-energy representations of $L G$ defined in Section 3. These representations are mutually locally unitarily equivalent, cf. [40] IV.6, [77] IV.2.4.1 and [84] p.12. Accordingly, for any $\lambda_i$ and any $I \in \mathcal{I}$ there is a unitary $u_I : \mathcal{H}_{\lambda_0} \to \mathcal{H}_{\lambda_i}$ such that $\lambda_i(g) = u_I\lambda_0(g)u_I^*$ for all $g \in L_I G$. Then, we can define an irreducible representation $\pi_{\lambda_i}$ of $\mathcal{A}_{G_\ell}$ by
\begin{equation}
\pi_{\lambda_i,I}(x) := u_I x u_I^*, \quad x \in \mathcal{A}_{G_\ell}(I), I \in \mathcal{I}.
\end{equation}

Note that it follows from Lemma 4.1 that, for any $I \in \mathcal{I}$, $\pi_{\lambda_i,I}$ does not depend on the choice of $u_I$ so that $\pi_{\lambda_i,I_2}|_{\mathcal{A}_{G_\ell}(I_1)} = \pi_{\lambda_i,I_1}$ whenever $I_1 \subset I_2$.

Henceforth, we make the following standing assumption on $\mathcal{A}_{G_\ell}$:

**Assumption 4.2.** $\mathcal{A}_{G_\ell}$ is completely rational and there exists a (necessarily unique) ring isomorphism $\psi_{G_\ell} : \mathbb{R}(L G) \to \hat{\mathbb{R}}\mathcal{A}_{G_\ell}$ such that $\psi_{G_\ell}([\lambda_0]) = [\pi_{\lambda_0}], i = 0, \ldots, N-1$.

This assumption might seem a strong restriction, but it is expected to be true in general and in all explicitly computed cases it has been proven [50] Sect.3.2, see also [57] Problem 3.32. E.g. for $G = \text{SU}(n)$ at any level $\ell$ the assumption follows from the results in [61, 86, 90].

Coming back to the case of general $G$ and $\ell$ with the above assumptions, we have the following immediate consequence of [12] Sect.3:

**Proposition 4.3.** If $\mathcal{A}_{G_\ell}$ satisfies Assumption 4.2 then the following holds:

(i) $C^\text{red}(\mathcal{A}_{G_\ell}) = B_{G_\ell}$ and $\mathfrak{R}_{\mathcal{A}_{G_\ell}} = \hat{\mathfrak{R}}_{G_\ell}$, so that $\mathfrak{R}_{\mathcal{A}_{G_\ell}} \supset \hat{\mathfrak{R}}_{G_\ell}$.

(ii) Every localized endomorphism $\rho$ of $C^*(\mathcal{A}_{G_\ell})$ with finite statistical dimension gives rise to a localized endomorphism $\hat{\rho}$ of $B_{G_\ell}$, which restricts to an endomorphism of the $C^*$-subalgebra of compacts $\hat{\mathfrak{R}}_{G_\ell} \subset B_{G_\ell}$.

Recall the injective ring homomorphism $\phi_0 : \mathbb{R}(L G) \to KK(\hat{\mathfrak{R}}_{G_\ell}, \hat{\mathfrak{R}}_{G_\ell})$ defined by (3.18). We now state and prove the main result of the present section. It gives an interpretation of $\phi_0$ in terms of DHR endomorphisms and, thanks to Remark 3.9 it also gives a description of $\phi_{-1}$ and $\phi_1$ in terms of DHR endomorphisms.

**Theorem 4.4.** Let $\mathcal{A}_{G_\ell}$ satisfy Assumption 4.2. Then, for every localized covariant endomorphism $\rho$ of $C^*(\mathcal{A}_{G_\ell})$ with finite statistical dimension we have $\phi_0(\psi_{G_\ell}^{-1}([\pi_0 \circ \rho])) = \{\hat{\rho} |_{\hat{\mathfrak{R}}_{G_\ell}}\}$.

**Proof.** For every $i = 0, \ldots, N-1$ choose a covariant localized endomorphism $\rho_i$ of $C^*(\mathcal{A}_{G_\ell})$ such that $[\pi_0 \circ \rho_i] = [\pi_{\lambda_i}]$. By the proof of [11] Thm.3.1 there is a surjective ring isomorphism $j : \hat{\mathfrak{R}}_{\mathcal{A}_{G_\ell}} \to KK(\hat{\mathfrak{R}}_{G_\ell}, \hat{\mathfrak{R}}_{G_\ell})$ such that $j([\pi_0 \circ \rho]) = \{\hat{\rho} |_{\hat{\mathfrak{R}}_{G_\ell}}\}$ for every covariant localized endomorphism $\rho$ with finite statistical dimension. Accordingly, we only need to show that $\phi_0([\lambda_i]) = \{\hat{\rho}_i |_{\hat{\mathfrak{R}}_{G_\ell}}\}$ for all $i \in \{0, \ldots, N-1\}$. Moreover, thanks to the isomorphism $KK(\hat{\mathfrak{R}}_{G_\ell}, \hat{\mathfrak{R}}_{G_\ell}) \simeq \text{End}(K_0(\hat{\mathfrak{R}}_{G_\ell}))$ it is enough to show that $[p_{\lambda_i}] \times \phi_0([\lambda_i]) = [p_{\lambda_i}] \times \{\hat{\rho}_i |_{\hat{\mathfrak{R}}_{G_\ell}}\}$, for all $i, j \in \{0, \ldots, N-1\}$. From [11] Thm.2.1(7) we see that $[p_{\lambda_i}] \times \{\hat{\rho}_i |_{\hat{\mathfrak{R}}_{G_\ell}}\}$ is equal to $[\hat{\rho}_i(\lambda_j)]$. Moreover, it follows from (4) in the
proof of [12] Thm.4.4] that \( [p_{\lambda_j}] = [\hat{\rho}_T(p_{\lambda_0})] \) and that
\[
[\hat{\rho}_i(p_{\lambda_j})] = [\hat{\rho}_i \circ \hat{\rho}_T(p_{\lambda_0})] = \sum_{k=0}^{N-1} \mathcal{N}^k_{i,j}[\hat{\rho}_k(p_{\lambda_0})] = \sum_{k=0}^{N-1} \mathcal{N}^k_{i,j}[p_{\lambda_0}] = [p_{\lambda_j}] \times \phi_0([\lambda_i]),
\]
where, for the last equality, we made use of (3.19). □

**Remark 4.5.** Every *-endomorphism \( \beta \) of \( \mathcal{R}_{G_i} \) induces an endomorphism \( \beta_* \) of \( K_0(\mathcal{R}_{G_i}) \) as push-forward, defined by \( \beta_*([p]) := [\beta(p)] \); it also induces an endomorphism \( \beta^* \) of \( K^0(\mathcal{R}_{G_i}) \) as pull-back: for a given K-homology class
\[
\varepsilon = \left[ \mathcal{H}, \pi \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \in K^0(\mathcal{R}_{G_i})
\]
it is defined by
\[
\beta^*(\varepsilon) := \left[ \mathcal{H}, \pi \circ \beta \oplus 0, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \in K^0(\mathcal{R}_{G_i}).
\]
According to [3] Ex.18.4.2(a)-(b)],
\[
[p] \times \{ \hat{\rho} \upharpoonright \mathcal{R}_{G_i} \} = (\hat{\rho} \upharpoonright \mathcal{R}_{G_i})_*([p]), \quad [p] \in K_0(\mathcal{R}_{G_i}),
\]
and
\[
\{ \hat{\rho} \upharpoonright \mathcal{R}_{G_i} \} \times \varepsilon = (\hat{\rho} \upharpoonright \mathcal{R}_{G_i})^*(\varepsilon), \quad \varepsilon \in K^0(\mathcal{R}_{G_i}).
\]
Recalling Remark 3.9 and Theorem 4.4, we can therefore express \( \phi_{-1} \) and \( \phi_1 \) as
\[
\phi_{-1}(\psi_{G_i}^{-1}([\pi_0 \circ \rho])) = (\hat{\rho} \upharpoonright \mathcal{R}_{G_i})_*([p_1]), \quad \phi_1(\psi_{G_i}^{-1}([\pi_0 \circ \rho])) = (\hat{\rho} \upharpoonright \mathcal{R}_{G_i})^*(\varepsilon_{\lambda_0}).
\]

**Remark 4.6.** In the setting of the preceding theorem, suppose that moreover \( \hat{\rho} \) preserves the differentiable subalgebra \( \mathcal{R}^1_{G_i} \subset \mathcal{R}_{G_i} \). Then it induces a pull-back endomorphism of \( HE^e(\mathcal{R}_{G_i}) \), which we denote by \( (\hat{\rho} \upharpoonright \mathcal{R}_{G_i})^* \). The two explicit cocycles \( (\hat{\rho}_i \upharpoonright \mathcal{R}_{G_i})^*\tau_{\lambda_0} \) and \( \tau_{\lambda_i} \) have the same pairing with \( K_0(\mathcal{R}_{G_i}) = K_0(\mathcal{R}^1_{G_i}) \), although they do not coincide. This construction was studied in [15] in a related context though with different underlying algebras.

**Remark 4.7.** If \( G \) is a connected simply connected compact Lie group which is not necessarily simple then \( G \) is the direct product of connected simply connected compact simple Lie groups and the results in this section and in Section 3 generalizes in a straightforward way. In this case the level \( \ell = (\ell_1, \ell_2, \ldots, \ell_n) \) consists of a level \( \ell_i \) (a positive integer) for each simple factor \( G_i \). Then the net \( \mathcal{A}_{G_i} \) is the tensor product \( \mathcal{A}_{G_1,\ell_1} \otimes \mathcal{A}_{G_2,\ell_2} \cdots \otimes \mathcal{A}_{G_n,\ell_n} \) where \( \mathcal{A}_{G_i,\ell_i}, i = 1, 2, \ldots, n \) denotes the net defined from the vacuum representation of the loop group \( LG_i \) at level \( \ell_i \). If \( \mathcal{A}_{G_i,\ell_i} \) satisfies Assumption 4.2 for \( i = 1, \ldots, n \), then also \( \mathcal{A}_{G_i} \) does.

5. **Non-simply connected compact Lie groups and other CFT models**

In this section we discuss the generalization of the results in Section 3 and Section 4 to other CFT models. The strategy is the following. We first give an abstract formulation of the results in terms of conformal nets admitting suitable supersymmetric extensions. We then show that the results apply to a large class of CFT models including lattice models, loop group models associated to non-simply connected compact Lie groups, coset models, the moonshine conformal net having the monster group \( M \) as automorphism group and the even shorter moonshine net having the baby monster group \( B \) as automorphism group.
5.1. Conformal nets with superconformal tensor product. In this subsection we will need the notion of graded-local conformal net (also called Fermi conformal net) which is a generalization of the one of conformal net when the axiom of locality is relaxed to graded-locality (or super-locality). These are the operator algebraic analogue of vertex operator superalgebras. We will need also to consider the spacial case of superconformal nets in which the conformal symmetry admits a supersymmetric extension. Basically this means that the vacuum Hilbert space of the net carries a representation of the Neveu-Schwarz super-Virasoro algebra compatible with the diffeomorphism symmetry of the net. For the precise definitions we refer the reader to [15, 16], cf. also [14].

**Definition 5.1.** Let $A$ be a conformal net. We say that $A$ admits a superconformal tensor product if there is a graded-local net $B$ with a graded positive-energy Ramond representation $\pi^B$, cf. [13, Thm. 2.13], satisfying the following properties:

(i) The graded-local conformal net $A \otimes B$ is superconformal in the sense of [15, Definition 2.11].

(ii) The Ramond representation $\pi^B$ satisfies the trace-class condition i.e. $e^{-tL_0^B}$ is a trace-class operator for all $t > 0$.

We will say that $A \otimes B$, or more precisely the pair $(A \otimes B, \pi^B)$, is a superconformal tensor product for the (local) conformal net $A$.

**Remark 5.2.** Note that the superconformal tensor products for a given conformal net $A$ are far from being unique. In fact if $(A \otimes B, \pi^B)$ is a superconformal tensor product for $A$ and $C$ is any superconformal net with a graded Ramond representation $\pi^C$ satisfying the trace class condition, then $(A \otimes (B \otimes C), \pi^{B \otimes C})$, where $\pi^{B \otimes C}$ is any graded subrepresentation of $\pi^B \otimes \pi^C$, is again a superconformal tensor product for $A$. Here, $\otimes$ denotes the graded tensor product, see e.g. [16, Subsec.2.6]. In particular one could take $C = A \otimes B$ and $\pi^C = \pi_0 \otimes \pi^B$, where $\pi_0$ is the vacuum representation of $A$. Accordingly, if the net $A$ admits a superconformal tensor product then it admits infinitely many superconformal tensor products.

**Remark 5.3.** The notion of superconformal tensor product can be defined in a completely analogous way for vertex operator algebra operators.

The motivating examples for the above definition are the conformal nets $A_{G_t}$ considered in Section 3. Let $F$ be the graded local conformal net on $S^1$ generated by a real free Fermi field (the free Fermi net), see e.g. [15,16]. Then for every positive integer $n$, the net $F^n$ generated by $n$ real free Fermi fields can be defined inductively by $F^1 := F$ and $F^{n+1} := F^n \otimes F$, $n \in \mathbb{N}$. The super-Sugawara construction described Section 6 shows that the net $A_{G_t}$ admits a superconformal tensor product $(A_{G_t} \otimes B, \pi^B)$, where $B = A_{G_t} \otimes F^d$, with $d$ the dimension of $G$ and $\pi^B = \pi_0 \otimes \pi_R$ with $\pi_0$ the vacuum representation of $A_{G_t}$ and $\pi_R$ the irreducible Ramond representation of $F^d$. Another possible choice is $B = F^d$ and $\pi^B = \pi_R$, where $\pi_R$ is the unique irreducible representation of $F^d$ if $d$ is even or the direct sum of the two inequivalent irreducible Ramond representations of $F^d$ if $d$ is odd. Then, in this case, the superconformal nets $A_{G_t} \otimes F^d$ are the super-current algebra nets considered in [15, Sec.6]. Note that for $d$ odd the irreducible Ramond representations of $F^d$ are not graded and for this reason we have to chose a reducible Ramond representation $\pi_R$. However, the irreducibility of $\pi^B$ in a superconformal tensor product $(A \otimes B, \pi^B)$ is not necessary for the purposes of this paper. The important property is the trace-class condition for $\pi^B$.

Many other examples can be given thanks to the following two propositions whose proofs are rather straightforward and will be omitted here.
Proposition 5.4. Let $\mathcal{A}$ be a conformal net on $S^1$ with a superconformal tensor product $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$. If $\hat{\mathcal{A}}$ is an irreducible local extension of $\mathcal{A}$ then $(\hat{\mathcal{A}} \otimes \mathcal{B}, \pi^B)$ is a superconformal tensor product for $\hat{\mathcal{A}}$.

Proposition 5.5. If $\mathcal{A}_1, \mathcal{A}_2$ are conformal nets with superconformal tensor products $(\mathcal{A}_1 \otimes \mathcal{B}_1, \pi^B_{1})$ and $(\mathcal{A}_2 \otimes \mathcal{B}_2, \pi^B_{2})$ respectively then, for every graded subrepresentation $\pi^B_{1} \otimes \pi^B_{2}$ of $\pi^B_{1} \otimes \pi^B_{2}$, $(\mathcal{A}_1 \otimes \mathcal{A}_2) \otimes (\mathcal{B}_1 \otimes \mathcal{B}_2)$ is a superconformal tensor product for the local conformal net $\mathcal{A}_1 \otimes \mathcal{A}_2$ with $\pi^B_0$ the vacuum representation.

Now, let $\mathcal{A}$ be a completely rational conformal net with $N$ irreducible sectors and let $\{\pi_0, \pi_1, \ldots, \pi_{N-1}\}$ be a maximal family of irreducible locally normal representations of $\mathcal{A}$, with $\pi_0$ the vacuum representation. Moreover, for any $i = 0, 1, \ldots, N-1$ let $\rho_i$ be a localized covariant endomorphism of the universal $C^*$-algebra $\mathcal{C}^*(\mathcal{A})$ such that $[\pi_i] = [\pi_0 \circ \rho_i]$. The reduced representation $\pi_{\mathrm{red}}$ is the direct sum of $\pi_0, \pi_1, \ldots, \pi_{N-1}$ defined on the Hilbert space $H_{\pi_{\mathrm{red}}} := H_{\pi_0} \oplus H_{\pi_1} \oplus \cdots \oplus H_{\pi_{N-1}}$.

The compact universal $C^*$-algebra $\mathcal{K}_{\mathcal{A}}$ is then the subalgebra of the algebra of compact operators $K(H_{\pi_{\mathrm{red}}})$ given by the direct sums of the algebras $K(H_{\pi_i})$, $i = 0, \ldots, N-1$. For any $i = 0, \ldots, N-1$, we choose a lowest energy unit vector $\Omega_i$ in $H_{\pi_i}$, and denote by $q_{\Omega_i}$ the corresponding one-dimensional projection in $K(H_{\pi_i})$ and let $p_i$ be the unique minimal projection in $\mathcal{C}^*(\mathcal{A})$ such that $\pi_i(p_i) = q_{\Omega_i}$. By [11, 12], see also Section 3 in this paper, the maps $[\pi_i] \mapsto [p_i] \in K_0(\mathcal{K}_{\mathcal{A}})$ and $[\pi_i] \mapsto \{p_i|_{\mathcal{K}_{\mathcal{A}}]\} \in KK(\mathcal{K}_{\mathcal{A}}, \mathcal{K}_{\mathcal{A}})$ give rise to a (surjective) group isomorphism

$$\phi^A_i : \hat{\mathcal{K}_{\mathcal{A}}} \to K_0(\mathcal{K}_{\mathcal{A}})$$

and to an injective ring homomorphism

$$\phi^A_0 : \hat{\mathcal{K}_{\mathcal{A}}} \to KK(\mathcal{K}_{\mathcal{A}}, \mathcal{K}_{\mathcal{A}}).$$

Moreover, $\phi^A_i$ and $\phi^A_0$ are related through the Kasparov product by

$$\phi^A_i(x) = [p_i] \times \phi^A_0(\hat{x}) , \quad x \in \mathcal{K}_{\mathcal{A}}.$$  

In order to define a group isomorphism $\phi^A : \hat{\mathcal{K}_{\mathcal{A}}} \to K_0(\mathcal{K}_{\mathcal{A}})$ by means of Dirac operators as in Sections 3 and 4 we assume that the completely rational conformal net $\mathcal{A}$ admits a superconformal tensor product $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$. We also assume that $\mathcal{A}$ has a trace-class representation theory, namely that $e^{-tL^B_{\hat{\pi}}}$ is trace class for all $t > 0$ and all irreducible locally normal representations $\pi$ of $\mathcal{A}$. All known completely rational conformal nets have a trace-class representation theory. Note also that if $\mathcal{A}$ is a modular net in the sense of [59], then $\mathcal{A}$ has a trace-class representation theory. Now, for every locally normal representation $\pi$ of $\mathcal{A}$ with finite statistical dimension we consider the Ramond representation $\hat{\pi}$ of the superconformal net $\mathcal{A} \otimes \mathcal{B}$ defined by $\hat{\pi} := \pi \otimes \pi^B$. Then, by our assumptions, the conformal Hamiltonian

$$L^\pi_0 = L^\pi_0 \otimes 1_{\mathcal{H}^B} + 1_{\mathcal{H}^A} \otimes L^B_0$$

has non negative spectrum and $e^{-tL^\pi_{\hat{\pi}}}$ is trace-class for all $t > 0$. Now, by the proof of [15, Prop. 2.14], there is a unitary positive-energy representation of the Ramond super-Virasoro algebra on $\mathcal{H}_\pi$ (cf. [3, 12]) by operators $L^\pi_n, G^\pi_n, n, r \in \mathbb{Z}$, with central charge $c = c_A + c_B$, where $c_A$ and $c_B$ are the central charges of $\mathcal{A}$ and $\mathcal{B}$ respectively. Then, the Dirac operator $L^\pi := G^\pi_0$ satisfies $D^2 = L^\pi - \frac{c}{24} 1_{\mathcal{H}^A} \geq (h_B - \frac{c}{24}) 1_{\mathcal{H}^A}$, where $h_B$ is the lowest energy in the representation $\pi^B$. Since $h_B$ is an eigenvalue of $L^\pi_0$ we see that we always have $h_B - \frac{c_A + c_B}{24} \geq 0$ and we say that the superconformal tensor product $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$ for $\mathcal{A}$ has the strict positivity property if $h_B - \frac{c_A + c_B}{24} > 0$. In this
case the Dirac operator $D_x$ has trivial kernel for every locally normal representation $\pi$ of $\mathcal{A}$ with finite statistical dimension.

**Proposition 5.6.** If a conformal net $\mathcal{A}$ admits a superconformal tensor product then it also admits a superconformal tensor product with the strict positivity property.

**Proof.** Let $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$ be a superconformal tensor product for $\mathcal{A}$ and let $\mathcal{C} := \mathcal{A}_G \otimes \mathcal{F}^d$ be a super-current algebra net, where $G$ is a simply connected compact simple Lie group, $d$ is the dimension of $G$ and the level $\ell$ is a positive integer. Moreover, let $\pi^C := \pi_0 \otimes \pi_R$ with $\pi_R$ the graded Ramond representation of $\mathcal{F}^d$ as defined above. Then $\mathcal{C}$ is a superconformal net with central charge $c_C = \frac{d}{2} + \frac{d}{\ell - h \omega}$ and $\pi^C$ is a graded Ramond representation of $\mathcal{C}$ with lowest energy $h_C = \frac{d}{16}$. Accordingly

$$h_C - \frac{c_C}{24} = d \left( \frac{1}{24} - \frac{\ell}{24(\ell + h \omega)} \right) > 0.$$ 

Now, let $\mathcal{B} := \mathcal{B} \otimes \mathcal{C}$ and $\pi^B := \pi^B \otimes \pi^C$. Then, $\mathcal{B}$ has central charge $c_B = c_B + c_C$ and $\pi^B$ has lowest energy $h_B + h_C$ so that

$$h_B - \frac{c_A + c_B}{24} = h_B - \frac{c_A + c_B}{24} + h_C - \frac{c_C}{24} \geq h_C - \frac{c_C}{24} > 0.$$ 

Hence, $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$ is a superconformal tensor product for $\mathcal{A}$ with the strict positivity property. \qed

Now let $\mathcal{A}$ be a completely rational conformal net with $N$ irreducible sectors and let $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$ be a superconformal tensor product for $\mathcal{A}$ having the strict positivity property. We fix a projection $e_B \in \mathcal{B}(\mathcal{H}_{\pi B})$ onto any one-dimensional even lowest energy subspace of $\mathcal{H}_{\pi B}$. Let $\pi$ be a locally normal representation of $\mathcal{A}$ with finite statistical dimension and let $\pi'$ be the unique normal representation of the reduced universal C*-algebra $C_{\text{red}}(\mathcal{A})$ such that $\pi' \circ \pi_{\text{red}} = \pi$. We define a degenerate representation $\tilde{\pi}$ of $\mathcal{R}_A$ on $\mathcal{H}_\pi = \mathcal{H}_\pi \otimes \mathcal{H}_{\pi B}$ by $\tilde{\pi}(x) := \pi'(x) \otimes e_B$. In particular we can define pairwise unitarily inequivalent representations $\tilde{\pi}_i$, $i = 0, \ldots, N - 1$. Using the Dirac operators $D_{\pi_i}$, we can define the Fredholm modules $(\mathcal{H}_{\tilde{\pi}_i}, \tilde{\pi}_i, \text{sgn} \ D_{\pi_i})$ for the C*-algebra $\mathcal{R}_A$ and the corresponding K-homology classes $\varepsilon_i := [(\mathcal{H}_{\tilde{\pi}_i}, \tilde{\pi}_i, \text{sgn} \ D_{\pi_i})] \in K^0(\mathcal{R}_A)$. Then, as in Proposition 3.7, it can be shown that the Kasparov product with the K-theory classes $[p_i]$ gives

$$[p_i] \times \varepsilon_j = \delta_{i,j}, \quad i, j = 0, \ldots, N - 1. \quad (5.2)$$

Consequently, the K-homology classes $\varepsilon_i$ do not depend on the choice of the superconformal tensor product $(\mathcal{A} \otimes \mathcal{B}, \pi^B)$. Now, the maps $[\pi_i] \mapsto \varepsilon_i$ give rise to a unique (surjective) group isomorphism $\phi^A_1 : \mathcal{R}_A \to K^0(\mathcal{R}_A)$. As a consequence of Eq. (5.2) and of the results in [11, 12], cf. the proof of Theorem 4.4, we have the following theorem.

**Theorem 5.7.** Let $\mathcal{A}$ be a completely rational conformal net having $N$ irreducible sectors and admitting a superconformal tensor product. Then the group isomorphism $\phi^A_1 : \mathcal{R}_A \to K^0(\mathcal{R}_A)$ does not depend on the choice of the superconformal product. Moreover,

$$\phi^A_0(x) \times \phi^A_1(y) = \phi^A_1(xy), \quad x, y \in \mathcal{R}_A.$$ 

In particular,

$$\phi_{-1}(x_k) \times \phi_0(x_i) \times \phi_1(x_j) = N_{ij}^k,$$

where $x_i := [\pi_i], i = 0, \ldots, N - 1$, denotes the preferred basis of $\mathcal{R}_A$. 
Remark 5.8. If \( \mathcal{A} \) is a completely rational conformal net with \( N \) irreducible sectors which has trace-class representation theory and admits a superconformal tensor product then, following the arguments in Section 3, one can define a differentiable Banach algebra \( \mathcal{R}^1_\mathcal{A} \subset \mathcal{R}_\mathcal{A} \) which is a dense subalgebra of \( \mathcal{R}_\mathcal{A} \) closed under holomorphic functional calculus so that \( K_0(\mathcal{R}^1_\mathcal{A}) = K_0(\mathcal{R}_\mathcal{A}) \). Moreover, as in Theorem 3.10 one can define even \( \theta \)-summable spectral triples \((\mathcal{R}^1_\mathcal{A}, \pi_i, D_{\pi_i})\) whose JLO cocycles \( \tau_i \) gives the same index pairing as the K-homology classes \( \varepsilon_i \), namely \( \tau_i(p_j) = [\tau_i] \times \varepsilon_i = \delta_{ij}, \quad i, j = 0, \ldots, N - 1 \). Accordingly the entire cyclic cohomology classes \( [\tau_i] \) give complete noncommutative geometric invariants for the irreducible sectors of \( \mathcal{A} \). However, in contrast to the group isomorphism \( \phi^A : \mathcal{R}_\mathcal{A} \to K^0(\mathcal{R}_\mathcal{A}) \) the differential algebra \( \mathcal{R}^1_\mathcal{A} \) and hence the corresponding cyclic cohomology classes \( [\tau_i] \) may depend on the choice of the superconformal tensor product for \( \mathcal{A} \).

Remark 5.9. If \( \mathcal{A} \) is an arbitrary completely rational conformal net with sectors \( [\pi_i], \quad i = 0, \ldots, N - 1 \) then one can define elements \( \varepsilon_i \in K^0(\mathcal{R}_\mathcal{A}) \) determined by Eq. (5.2). Then the map \( [\pi_i] \mapsto \varepsilon_i \) defines a group isomorphism \( \phi^A : \mathcal{R}_\mathcal{A} \to K^0(\mathcal{R}_\mathcal{A}) \) such that \( \phi^A(0) \times \phi^A(y) = \phi^A(xy) \) as in Theorem 5.7. However, without assuming that \( \mathcal{A} \) admits a superconformal tensor product and a trace-class representation theory we loose the natural interpretation of \( \phi^A \) in terms of Dirac operators, JLO cocycles and superconformal symmetry.

Remark 5.10. Let \( \mathcal{A} \) be a completely rational conformal net. Then \( K^0(\mathcal{R}_\mathcal{A}) \) is a finitely generated free abelian group. Every ring structure on \( K^0(\mathcal{R}_\mathcal{A}) \) obtained by introducing a product \( \ast \) compatible with the group operation \( + \) gives rise to a group homomorphism \( \phi^* : K^0(\mathcal{R}_\mathcal{A}) \to KK(\mathcal{R}_\mathcal{A}, \mathcal{R}_\mathcal{A}) \) determined by the condition \( x \ast y = \phi^*(x) \times y, \quad x, y \in K^0(\mathcal{R}_\mathcal{A}) \). Conversely every group homomorphism \( \phi : K^0(\mathcal{R}_\mathcal{A}) \to KK(\mathcal{R}_\mathcal{A}, \mathcal{R}_\mathcal{A}) \) determines a ring structure on \( K^0(\mathcal{R}_\mathcal{A}) \) through the product \( x \ast \phi y = \phi(x) \times y \). With the special choice \( \phi := \phi^A \circ (\phi^A)^{-1} \), with \( \phi^A \) determined by the action of the DHR endomorphisms on \( \mathcal{R}_\mathcal{A} \), the group isomorphism \( \phi^A : \mathcal{R}_\mathcal{A} \to K^0(\mathcal{R}_\mathcal{A}) \) becomes a ring isomorphism.

5.2. Applications to chiral CFT models. In this subsection we give various examples of conformal nets admitting superconformal tensor products. For the examples for which the net is known to be completely rational we can apply Theorem 5.7 and Remark 5.8, so that the DHR fusion ring of the nets can be described in terms of K-theory and noncommutative geometry. All examples considered below, completely rational or not, have a vertex operator algebra analogue.

Example 5.11. Let \( \mathcal{A}_{U(1)} \) be the conformal net generated by a chiral U(1) current (chiral free Bose field) considered in [9], see also [17] Example 8.6. Then it follows from the super-Sugawara construction in [50], cf. Section 3, and also [15] Sec.6, that \((\mathcal{A}_{U(1)} \otimes \mathcal{F}, \pi_R)\) is a superconformal tensor product for \( \mathcal{A}_{U(1)} \), where \( \mathcal{F} \) is the free Fermi net and \( \pi_R \) is the corresponding graded Ramond representation. More generally, for any positive integer \( n \), the conformal net \( \mathcal{A}_{U(1)^n} \) defined as the tensor product of \( n \) copies of \( \mathcal{A}_{U(1)} \), admits the superconformal tensor product \((\mathcal{A}_{U(1)^n} \otimes \mathcal{F}^n, \pi_R)\). All the nets \( \mathcal{A}_{U(1)^n} \) admit uncountably many irreducible sectors, see e.g. [9], and hence cannot be completely rational.

Example 5.12. Let \( L \) be an even positive-definite lattice of rank \( n \). Then one can define a corresponding conformal net \( \mathcal{A}_L \) [26], [76], which is completely rational by [26] Corollary 3.19. The conformal net \( \mathcal{A}_L \) is the operator algebraic analogue of the simple lattice vertex operator algebra \( V_L \), see e.g. [53] Sect.5.5 for the definition of \( V_L \), see also [17] Conjecture 8.17]. By construction, the net \( \mathcal{A}_L \) is an irreducible local extension of \( \mathcal{A}_{U(1)^n} \) and hence it admits a superconformal tensor product. In fact one can choose \((\mathcal{A}_L \otimes \mathcal{F}^n, \pi_R)\).
Remark 5.13. The lattice net $\mathcal{A}_L$ and its representations are related to the projective unitary positive-energy representations of the loop group $L U(1)^n$ corresponding to a central extension determined by $L$ through the group isomorphism between $U(1)^n$ and the $n$-dimensional torus $\mathbb{R}L/L$, see [72] Sect.9.5 and [25] Sect.3. Accordingly $\mathcal{A}_L$ may be considered as a loop group net for the non-simply connected group $U(1)$. The subnet $\mathcal{A}_{U(1)^n} \subset \mathcal{A}_L$ then corresponds to the restriction to the component of the identity $(L U(1)^n)_1$ of $L U(1)^n$.

Example 5.14. Let $G$ be a compact connected Lie group. Then

$$G \simeq (G_1 \times G_2 \times \cdots \times G_m \times U(1)^n)/Z ,$$

where $G_i$, $i = 1, \ldots, m$ is a connected simply connected compact simple Lie group and $Z$ is a finite subgroup of the center $Z(G_1) \times Z(G_2) \times \cdots \times Z(G_m) \times U(1)^n$ of $G_1 \times G_2 \times \cdots \times G_m \times U(1)^n$. Of course, $G$ is semisimple if and only if $n = 0$ i.e., there is no torus factor. Now, let $BG$ be the classifying space of $G$. A class $\ell \in H^4(BG, \mathbb{Z})$ is called a level and transgresses to a central extension $L G^\ell$ of $L G$, [36] [50] [83]. We say that the level $\ell \in H^4(BG, \mathbb{Z})$ is positive if $L G^\ell$ admits irreducible positive-energy unitary representations i.e., if there exists an irreducible positive-energy projective unitary representations at level $\ell$ [79], see also Definition 1 and Definition 2 in [50]. If $\ell \in H^4(BG, \mathbb{Z})$ is a positive level then one can define a local conformal net $\mathcal{A}_{G_\ell}$ as a simple current extension

$$\mathcal{A}_{G_\ell} := (\mathcal{A}_{G_1, \ell_1} \otimes \mathcal{A}_{G_2, \ell_2} \otimes \cdots \otimes \mathcal{A}_{G_m, \ell_m} \otimes \mathcal{A}_L) \rtimes Z ,$$

where the positive integers (levels) $\ell_1, \ldots, \ell_m$ and the lattice $L$ are determined by $\ell \in H^4(BG, \mathbb{Z})$, see [50], [71]. Equivalently, the net $\mathcal{A}_{G_\ell}$ can be defined through the vacuum representation $\lambda_0$ of $L G^\ell$ as in the simply connected case, see Eq. (1.2). The locality property of the net follows from the disjoint-commutativity of the “transgressive” central extension $L G^\ell$, see [83] Sect.3.3. Now, $\mathcal{A}_{G_{i, \ell_i}}$, $i = 1, \ldots, m$ and $\mathcal{A}_L$ admit a superconformal tensor product and hence, by Proposition 5.5 the net

$$\tilde{\mathcal{A}} := \mathcal{A}_{G_1, \ell_1} \otimes \mathcal{A}_{G_2, \ell_2} \otimes \cdots \otimes \mathcal{A}_{G_m, \ell_m} \otimes \mathcal{A}_L$$

admits a superconformal tensor product. Accordingly, since $\mathcal{A}_{G_{i, \ell_i}} = \tilde{\mathcal{A}} \rtimes Z$ is an irreducible local extension of $\tilde{\mathcal{A}}$, it also admits a superconformal tensor product by Proposition 5.4. Actually, the supersymmetric tensor product can be taken of the form $(\mathcal{A}_{G_{i, \ell_i}} \otimes \mathcal{F}^d, \pi_R)$ where $d$ is the dimension of $G$, $\mathcal{F}^d$ is the graded-local conformal net generated by $d$ real free Fermi fields and $\pi_R$ is a graded Ramond representation of $\mathcal{F}^d$. Hence, if $\mathcal{A}_{G_{i, \ell_i}}$ is completely rational, equivalently if $\mathcal{A}_{G_{i, \ell_i}}$ is completely rational for all $i = 1, \ldots, m$, Theorem 5.7 applies and gives a generalization of the results in Section 4 to the case of non-simply connected compact Lie groups.

Remark 5.15. Similarly to Example 5.14 one can define a vertex operator algebra $V_{G_{i, \ell_i}}$, which is the analogue of the conformal net $\mathcal{A}_{G_{i, \ell_i}}$, as a simple current extension of the tensor product of an affine vertex operator algebra and a lattice vertex operator algebra, cf. [25] [50] [51] [66].

Remark 5.16. Let $G$ and $\ell \in H^4(BG, \mathbb{Z})$ as in Example 5.14 and let $R^\ell(L G)$ be the free abelian group generated by the irreducible positive-energy unitary representations of $L G^\ell$, i.e., the level $\ell$ irreducible positive-energy projective unitary representations of $L G$, with $e^{2\pi i L_3^\ell}$ a multiple of the identity, see e.g. [36]. Note that the condition that $e^{2\pi i L_3^\ell}$ a multiple of the identity is not necessarily satisfied for general positive-energy representations, see [79]. Then $R^\ell(L G)$ admits a fusion ring structure which can be defined through the correspondence with the VOA simple modules of the vertex operator algebra $V_{G_i}$ or through the modular invariance property of characters as in the simply
connected case. Moreover, as in the simply connected case, any irreducible positive-energy unitary representation $\lambda$ of $LG^\ell$ is locally unitarily equivalent to the vacuum representation $\lambda_0$ and hence gives rise to an irreducible locally normal representation $\pi_\lambda$ of the conformal net $\mathcal{A}_{G\ell}$ as in Eq. (1.5) and, if $\mathcal{A}_{G\ell}$ is completely rational, the map $\lambda \mapsto \pi_\lambda$ gives rise to an injective group isomorphism $\psi_{G\ell}: R^d(LG) \to \widetilde{R}_{\mathcal{A}_{G\ell}}$. It is expected that $\mathcal{A}_{G\ell}$ is always completely rational and that $\psi_{G\ell}$ is always a surjective ring isomorphism, i.e. that $\mathcal{A}_{G\ell}$ satisfies the analogue of Assumption [12] for every connected compact Lie group $G$ and every positive level $\ell \in H^4(BG, \mathbb{Z})$, see e.g. [50, Conjecture 4], although this remains in general an important open problem. A positive solution is known for some special cases, e.g. for $G = SO(3) = SU(2)/\mathbb{Z}_2$ at every positive level, cf. [50, Sect. 3] and [5, Sect. 5].

**Example 5.17.** Let $G$ and $\ell \in H^4(BG, \mathbb{Z})$ be as in Example 5.14 and let $H \subset G$ be a closed connected subgroup of $G$. Then the positive level $\ell \in H^4(BG, \mathbb{Z})$ maps to a positive level $\ell' \in H^4(BH, \mathbb{Z})$ which gives a smooth central extension $LH^\ell \subset LG^\ell$ of $LH$. The restriction to $LH^\ell$ of the vacuum representation $\lambda_0$ of $LG^\ell$ gives rise to an embedding $\mathcal{A}_{H\ell'} \subset \mathcal{A}_{G\ell}$ of the conformal net $\mathcal{A}_{H\ell'}$ as a covariant subnet of $\mathcal{A}_{G\ell}$. The corresponding coset subnet $\mathcal{A}_{H\ell'}^G(I) = \mathcal{A}_{H\ell'}(I)' \cap \mathcal{A}_{G\ell}(I)$, $I \in \mathcal{I}$, see [92], see also [17, 68]. It follows from the Kazama-Suzuki superconformal coset construction [63, 64] that the coset conformal net $\mathcal{A}_{H\ell'}^G$ admits a superconformal tensor product $(\mathcal{A}_{H\ell'}^c \otimes \mathcal{F}^{d_G-d_H}, \pi_R)$, where $d_G$ and $d_H$ are the dimensions of $G$ and $H$ respectively. In various cases the coset net $\mathcal{A}_{H\ell'}^G$ is known to be completely rational, see [68, 91, 92, 93, 94]. In all these cases Theorem 5.7 applies. Moreover, the completely rational coset conformal nets admit interesting irreducible extensions such as the mirror extensions defined in [95]. Then all these extensions are completely rational and admit a superconformal tensor product so that Theorem 5.7 applies. Many interesting CFT models can be described through completely rational coset conformal nets and their irreducible extensions. Various examples will be given here below.

**Example 5.18.** Let $\mathcal{A}_{Vir,c}$ be the Virasoro net with central charge $c$ [10, 58]. If $\mathcal{A}$ is a conformal net then the corresponding representation of $Diff(S^1)$ gives rise to an irreducible subnet $\mathcal{A}_{Vir,c} \subset \mathcal{A}$, the Virasoro subnet of $\mathcal{A}$. The value $c$ is determined by $\mathcal{A}$ and $c$ is said to be the central charge of $\mathcal{A}$. If $c < 1$ then, as a consequence of the Goddard-Kent-Olive construction [44], $\mathcal{A}_{Vir,c}$ can be realized as a coset for an appropriate inclusion of loop group nets [58] and it turns out to be completely rational. Accordingly Theorem 5.7 applies to all Virasoro nets with $c < 1$ and, in fact, to all conformal nets with $c < 1$. These are classified in [58].

**Example 5.19.** The even (Bose) subnet of a $N = 1$ super-Virasoro net with central charge $c < 3/2$ can be realized as a coset of an inclusion of loop group nets and it turns out to be completely rational [10, Sect. 6]. Accordingly, the even subnets of the superconformal nets with $c < 3/2$ are all completely rational and admits a superconformal tensor product so that Theorem 5.7 applies. These conformal nets have been classified in [10, Sect. 7]. Similarly, the even subnet of a $N = 2$ super-Virasoro net with central charge $c < 3$ can be realized as a coset of an inclusion of loop group nets and it turns out to be completely rational [14, Sect. 5]. Accordingly, the even subnets of the $N = 2$ superconformal nets with $c < 3$ are all completely rational and admits a superconformal tensor product so that Theorem 5.7 applies. These conformal nets have been classified in [14, Sect. 6].

**Example 5.20.** For every subfactor $N \subset M$ with Jones index $[M : N] < 4$, M. Bischoff has constructed in [51] a completely rational conformal net $\mathcal{A}_{N\subset M}$ whose representation...
category is braided tensor equivalent to the quantum double $D(N \subset M)$ and has shown the existence of vertex operator algebras $V_{N \subset M}$ with the analogous property. The nets $A_{N \subset M}$ are obtained from loop group nets by taking cosets, irreducible local extensions and tensor products. Accordingly, they all admit a superconformal tensor product and Theorem $5.7$ applies. Note that, for any of these nets, the DHR fusion ring $\tilde{R}_{A_{N \subset M}}$ coincides with the fusion ring of the corresponding vertex operator algebra $V_{N \subset M}$.

**Example 5.21.** Let $n$ be a positive integer and let $A_{Vir,\frac{n}{2}}^\otimes$ be the conformal net with central charge $n/2$ defined as the tensor product of $n$ copies of the Virasoro net $A_{Vir,\frac{1}{2}}$. Then, $A_{Vir,\frac{n}{2}}^\otimes$ is the tensor product of completely rational conformal nets admitting a superconformal tensor product and hence it is a completely rational conformal net admitting a superconformal tensor product. A conformal net $A$ is said to be framed if it is an irreducible local extension of $A_{Vir,\frac{n}{2}}^\otimes$ for some positive integer $n$ [60 Sect.4], see also [62]. Accordingly, every framed conformal net is completely rational and admits a superconformal tensor product so that Theorem $5.7$ applies. Remarkable examples of framed conformal nets are the moonshine net $A_{Vir,\frac{n}{2}}^\otimes$ constructed by Kawahigashi and Longo in [60], see also [17 Thm.8.15] and whose automorphism group is the monster group $M$, and the even shorter net $A_{Vir,\frac{n}{2}}^\otimes$ constructed in [17 Thm.8.16].

All the above example show that the description of the representation theory of CFTs in terms of K-theory and noncommutative geometry goes far beyond the realm of loop groups. It would be very interesting if some of the above examples also admitted a topological description in terms of twisted K-theory in the spirit of FHT. This would give e.g. a twisted K-theory description of the discrete series representations of $\text{Diff}(S^1)$ and of the representation theory of coset models, cf. [29 page 2013] and [30 page 323]. To this end a more direct and clear relation of the results in this paper in the case of loop groups and the FHT work is probably needed.

We end this section with a comment on the case of disconnected compact Lie groups, which are covered in the FHT setting through the analysis of twisted loop groups and their positive-energy representations, cf. [30], but not in our present analysis. The point is that the usual definition of loop group nets generalizes to non-simply connected compact Lie groups but apparently does not generalize in a natural way to disconnected compact Lie groups, cf. [50] [51] [52] [53]. Note that in [52] [53] a fusion product on the representation of certain twisted loop group is constructed from the point of view of subfactor theory through Connes fusions. The results there indicate that there is no conformal net associated to a twisted loop group. Rather the representations of the twisted loop group should be considered as twisted (soliton) representations of the corresponding untwisted loop group net and should be related to the representation theory of orbifold models which, presently, are not covered by our analysis.

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