Lead palladium titanate: A room temperature nanoscale multiferroic thin film

K. K. Mishra1*, Alvaro A. Instan1, Shalini Kumari2, J. F. Scott3 & Ram S. Katiyar1*

The discovery of single-phase multiferroic materials and the understanding of coupling mechanisms between their spin and polarization is important from the point of view of next generation logic and memory devices. Herein we report the fabrication, dielectric, ferroelectric, piezo-response force microscopy, and magnetization measurements of Pd-substituted room-temperature magnetoelectric multiferroic PbPd0.3Ti0.7O3 (PbPdT) thin films. Highly oriented PbPdT thin films were deposited on [(LaAlO)0.3(SrAlTaO)0.7] (LSAT) substrates in oxygen atmosphere using pulsed laser deposition technique. X-ray diffraction studies revealed that the films had tetragonal phase with (001) orientation. Surface morphology studies using atomic force and scanning electron microscopy suggest a smooth and homogeneous distribution of grains on the film surface with roughness ~2 nm. A large dielectric constant of ~1700 and a low-loss tangent value of ~0.3 at 10 kHz were obtained at room temperature. Temperature dependent dielectric measurements carried out on Pt/PbPdT/La0.7Sr0.3MnO3 (LSMO) metal-dielectric-metal capacitors suggest a ferroelectric to paraelectric transition above 670 K. The measured polarization hysteresis loops at room temperature were attributed to its ferroelectric behavior. From a Tauc plot of (αhν)² versus energy, the direct band gap E_g of PbPdT thin films was calculated as 3 eV. Ferroelectric piezoelectric nature of the films was confirmed from a strong domain switching response revealed from piezo-response force microscopy. A well-saturated magnetization M-H loop with remanent magnetization of 3.5 emu/cm³ was observed at room temperature, and it retains ferromagnetic ordering in the temperature range 5–395 K. Origin of the magnetization could be traced to the mixed oxidation states of Pd²⁺/Pd⁴⁺ dispersed in polar PbTiO₃ matrix, as revealed by our x-ray photoelectron spectroscopic results. These results suggest that PbPdT thin films are multiferroic (ferroelectric-ferromagnetic) at room temperature.

Multiferroic materials with coupled electric and magnetic ordering in a single phase are of research interest due to their fascinating physics and potential for multifunctional device applications, such as speed writing and non-destructive data storage1–4. Only a limited number of single phase multiferroics are available due to the chemical incompatibility between magnetic and ferroelectric order parameters5. The most well-known examples are: BiFeO₃, Pb(Fe0.5Nb0.5)O₃, YMnO₃, TbMnO₃, GaFeO₃6–11. However, a majority of these materials exhibit multiferroicity at cryogenic temperature. Materials exhibiting multiferroicity at room temperature are important for device applications, but only a few have been discovered so far, the most celebrated examples being BiFeO₃10,12, epitaxial thin films of GaFeO₃13 and Bi₂FeCrO₆14–16. During last decade, extensive studies have been carried out to investigate BiFeO₃10–12,17 that exhibits both ferroics ordering above room temperature with ferroelectric to paraelectric phase transition temperature Tc ~ 1143 K and antiferromagnetic to paramagnetic phase transition Néel temperature T_N ~ 643 K, to understand the origin and mechanism of its multiferroicity. Pd-substituted PbTiO₃-based oxides are new promising room temperature multiferroic materials due to coexistence of ferroelectric and magnetization ordering at and above room temperature. Lead titanate (PbTiO₃) is a classical ferroelectric at room temperature, belonging to the perovskite family with the general formula ABX₃ (Fig. 1). The sublattices A is occupied by non-magnetic Pb²⁺ (6P⁰, S = 0) and the B-site is occupied by non-magnetic Ti⁴⁺ (3d⁰, S = 0) ions. It undergoes a first-order displacive phase transition from ambient tetragonal to high temperature cubic phase

1Department of Physics and Institute for Functional Nanomaterials, P.O. Box 70377, University of Puerto Rico, San Juan, PR, 00936-8377, USA. 2Department of Physics, The Pennsylvania State University, University Park, Pennsylvania, 16802, USA. 3School of Physics, University of St. Andrews, St. Andrews, KY19 9SS, UK. *email: karuna.kara@upr.edu; ram.katiyar@upr.edu
is oversized for Pd cations. Earlier studies suggested that Pd atom is as such non-magnetic (4d10, S = 0)26, however, under the influence of external electric field and/or strain, Pd cation can become ferromagnetic and its magnetic ordering can be electrically tunable21–24. A single phase of 30% Pd-substituted Pb(Zr0.20Ti0.80)0.70Pd0.30O3 ceramic was successfully prepared by solid state reaction method, that is reported to be room temperature multiferroic14. It exhibited a weak ferromagnetism, improved ferroelectricity, and importantly showed strong magneto-electric (ME) coupling. Very recently, Pd-doping of up to 30% in ferroelectric PbTiO3 was realized experimentally that established its magnetoelectric multiferroicity at room temperature, superior to BiFeO3.25. The magnetism in the ceramic sample was ascertained from the presence of Pd4+ chemical state from x-ray photo electron spectroscopy (XPS). However, thin films fabrication of this intriguing multiferroic and investigation of their polarization and magnetization (multiferroic order parameters) are yet to be carried out. Multifunctional thin films are preferable for device applications such as in highly sensitive actuators and sensors, multistate memories, and nanoelectronics10,26–29, and they are also fascinating for their rich physics. Pulsed laser deposition technique is a sophisticated thin film growth technique often useful for the fabrication of high purity thin films with nanoscale precision20–23 and is preferable for deposition of high quality films including epitaxial films on appropriate substrates.

We report the fabrication of PbPd0.3Ti0.7O3 (PbPdT) thin films using pulsed laser deposition technique and investigations of their structural, dielectric, ferroelectric, and magnetic ordering behavior with temperature. We have fabricated Pt/PbPdT/La0.7Sr0.3MnO3 (LSMO)/{(LaAlO3)0.3(Sr2AlTaO6)0.7} (LSAT) heterostructures for electrical measurements, and studied the magnetization of PbPdT thin films (of similar thickness) grown on LSAT (001) substrates in order to avoid any magnetization contribution from La0.7Sr0.3MnO3 (LSMO). These films had c-axis orientation. Study of the role of chemical substitution of Pd ions on the magnetic properties of the films revealed that the PbPdT films exhibited well defined M-H hysteresis loops up to 395 K and retained ferromagnetic ordering from 5 to 395 K, the highest temperature in the present study. X-ray photo electron spectroscopic measurements revealed mixed oxidation states of Pd2+ and Pd4+ in Pd-substituted polar PbTiO3 matrix and fulfilled the requirement to realize ferromagnetism in PbPdT films. Dielectric studies on Pt/PbPdT/LSMO metal-insulator-metal (MIM) capacitors suggest that the ferroelectric to paraelectric phase transition temperature is above 670 K. From the Tauc plot analysis of absorption spectra using (αhν)2 versus energy, the direct band gap $E_g$ of PbPdT thin films is estimated as 3 eV. The ferroelectric piezoelectric nature of the films is evident from the strong domain switching response obtained from the piezoresponse force microscopy (PFM) studies involving color contrast phase and amplitude images. Furthermore, magnetocapacitance effect was noticed in our thin films.

**Figure 1.** ABO3 type structure of PbTiO3: drawn using VESTA software47.

Results and Discussion

The PLD deposition parameters were optimized for the growth of PbPdT thin films as summarized in Table 1. The high-resolution x-ray diffraction pattern of one of the PbPdT thin films deposited on LSMO buffer layer coated on LSAT substrate is shown in Fig. 2. As can be seen, in the $\theta$-2$\theta$ scan, the reflections corresponding to PbPdT, LSMO and LSAT substrate were noticed in the 2$\theta$ range from 20°–60°. In addition, absence of any impurity peaks originating from other secondary phases indicated that the grown PbPdT thin film is phase pure and stabilized in a single perovskite phase. The thin films were stabilized in the tetragonal phase ($P4mm$), as in pristine PbTiO3 (JCPDF # 742495), and these were oriented with their (001) plane parallel to the substrate surface along with their out-of-plane lattice parameter $c$ as 3.958 Å. Lattice mismatch between the thin film and the substrate induces a structural strain, which can be estimated using the relation $\eta = (a_{\text{substrate}} - a_{\text{film}})/a_{\text{film}} \times 100$. The LSAT and LSMO stabilized in cubic perovskite phase with their lattice parameters $a_{\text{substrate}} = 3.868 \text{ Å}$ and $a_{\text{LSMO}} = 3.871 \text{ Å}$, respectively. However, the in-plane lattice parameters ($a$ and $b$) of the films cannot be extracted from the present...
In order to obtain these lattice parameters, in-plane x-ray diffraction and/or Reciprocal Space Mapping measurements are required. In the absence of these extra experimental results, the in-plane lattice strain in the films could not be calculated. The thickness of the films, estimated using a Profilometer, was found to be ~300 nm. The presence of the constituent elements, such as Pb, Pd, Ti, and O in the films were inferred from an energy dispersive x-ray spectrum (EDS) (Fig. 3) excited using an electron beam of kinetic energy 20 keV. X-ray emission lines corresponding to O Kα, 0.53 keV, Pb Mα, 2.34 keV, Pd Lα, 2.88 keV, Pd Lβ, 3.01 keV, Ti Kα, 4.53 keV and Pb Lα, 10.58 keV exactly match with the available theoretical characteristic x-ray emission lines. The surface topography of the thin films was measured using atomic force microscopy (AFM) (Fig. 3: inset) from the surface of the thin films indicating a smooth surface of the films with its root mean square roughness (Rq) of ~2 nm. The surface morphology of the films using scanning electron microscopy (SEM) is shown in Fig. 4a. It suggests that the thin films are smooth and dense with close packed crystal grains. A homogeneous distribution of these chemical elements was ascertained via SEM-based elemental mapping profiles (Fig. 4b–f), scanned over a large surface area of the films.

As mentioned earlier, it was argued from DFT studies as well as other earlier experimental reports on Pd-based ceramics that the origin of magnetism in PbPdT is due to the existence of Pd⁴⁺ cationic state at Pb²⁺ site (A-site). Therefore, it is imperative to examine the valence state of palladium (Pd) using x-ray photoelectron spectroscopy. From the survey spectra, the presence of all constituent elements, as inferred from EDS spectra, were confirmed from their respective binding energy matching with the standard theoretical values. The high resolution XPS spectrum of Pd consists of two components due to spin-orbit coupling effect (Fig. 5). The spin-orbit doublets of Pd were found to be asymmetric with peak centers at Pd 3d⁵/₂ = 335.65 eV and Pd 3d³/₂ = 340.91 eV, suggesting the presence of Pd⁴⁺ and Pd²⁺ valence states. These doublets were deconvoluted using the Casa XPS software and by choosing a suitable background. The deconvolution of Pd 3d⁵/₂ peak yielded Pd²⁺ state at 335.56 eV, and Pd⁴⁺ state at 341.37 eV (Fig. 5). The effective area ratio Pd⁴⁺:Pd²⁺ was calculated to be 0.83: 1. The smaller value of Pd⁴⁺ states can be explained by oxygen loss from the surface, that turns some Pd⁴⁺ into Pd²⁺ states.

Table 1. Growth parameters for PbPd₀.₃Ti₀.₇O₃ thin films deposited using pulsed laser deposition technique.

| Substrate | [(LaAlO₃)₀.₃(Sr₂AlTaO₆)₀.₇] (001) |
|-----------|--------------------------------|
| Target    | Pb₀.₇Pd₀.₃Ti₀.₇O₃               |
| Target diameter | 0.025 m |
| Substrate target distance | 0.05 m |
| Growth temperature | 923 K |
| Base vacuum | ~10⁻⁶ Torr |
| O₂ partial pressure | 200 mTorr |
| Laser source | excimer gas (λ = 248 nm, KrF) |
| Laser energy used | 250 mJ |
| Pulse repetition rate | 10 Hz |
| Total number of shots | 3000 |
| Post deposition ex-situ annealing | 923 K for 0.5 hour in O₂ atmosphere |

Figure 2. X-ray diffractogram measured in θ-2θ geometry for a PbPd₀.₃Ti₀.₇O₃ film grown on LSMO coated LSAT (001) substrate. The reflection peaks are also indexed for PbPd₀.₃Ti₀.₇O₃.
The observed magnetism in our PbPdT thin films is due to mixed oxidation states of Pd that essentially fulfills the requirement to realize the ferromagnetism in Pd-based perovskites. The dielectric constant ($\varepsilon$) and the loss tangent (tan $\delta$) of PbPdT thin films were measured in the frequency range 10$^2$–10$^6$ Hz and temperature between 100–670 K (Fig. 6(a)). The room temperature dielectric constant and dielectric loss were found to be ~1700 and ~0.3, respectively, measured at 10 kHz. The dielectric constant $\varepsilon$ was found to be higher than the reported values for PZTP (~300) and PbPdT sintered bulk samples (~500). One can notice a nominal decrease in dielectric constant and almost constant loss tangent values with increasing frequency up to 10$^5$ Hz. However, a substantial decrease in the dielectric constant and a rise in dielectric loss were observed above 10$^5$ Hz frequency. The reduction in $\varepsilon$ is expected since the dielectric constant of PbPdT thin films has contributions from intrinsic factors related to the lattice, and extrinsic contributions involving grain boundary and interface. In particular, the intrinsic contribution is dependent upon the grain size, film orientation, and strain on the films. However, the interface between the PbPdT film and the bottom LSMO electrode acts as a pinning center that contributes extrinsically to the dielectric constant. It perturbs the domain wall motion, and consequently leads to a reduction in the dielectric constant. Such a reduction of dielectric constant at high frequencies was reported in Pb(Zr$_{0.52}$Ti$_{0.48}$)O$_3$/LSMO thin films, and it was argued to arise due to interfacial contributions. A close inspection of the dielectric behavior of the reported LSMO films, and the present PbPdT/LSMO films at high frequencies indicate a similar decreasing trend suggesting that the dielectric response of our thin film capacitors (~10$^5$ Hz) could be mainly from the bottom layer (LSMO). The loss tangent values become larger at higher frequency which can be attributed to the semiconducting nature of the bottom layer, and also to the presence of disorder or defects in the films. The present dielectric behavior is similar to those reported in Pb(Zr$_{0.52}$Ti$_{0.48}$)O$_3$/LSMO thin films.
in other perovskite ferroelectrics, such as Pb(Zr_{0.53}Ti_{0.47})_{0.9}Sc_{0.1}O_{3} \cite{27}, Pb(Zr_{0.2}Ti_{0.8})_{0.3}O_{3-\delta} \cite{36}. At high temperatures (\geq 400 K), a substantial drop in the dielectric constant at high frequencies (>10^5 Hz) is observed due to thermal activation of hopping motion of oxygen vacancies. As in several perovskite oxides\cite{37,39}, ionic conductivity in PbPdT films is expected to increase due to increase in movement of oxygen vacancies with thermal energy. The temperature dependent behavior of dielectric constant and loss tangent in the frequency range 100 Hz - 1 MHz is shown in Fig. 6(b). The dielectric constant was found to be almost constant up to 400 K, followed by a drastic increase above this temperature. A sharp rise in dielectric constant above 670 K could indicate that the system is approaching the ferroelectric-paraelectric phase transition temperature. Using our temperature controller MMR K-20, the heater power allows us to reach only up to 670 K, therefore the actual dielectric maximum that is likely to be above this temperature (\geq 670 K), could not be observed. Indeed the ferroelectric Curie temperature T_c of PbPdT ceramics is reported to be 740 K \cite{25}. Pd cation is larger in size compared to Ti cation, therefore upon substituting such a large size Pd cation at the Ti-site in PbTiO_3, the unit cell volume increases. Hence the transition temperature T_c is expected to reduce compared to the pristine PbTiO_3 films (T_c \sim 823 K) \cite{40}. The rise of loss tangent at higher temperatures is due to space charge polarization contribution.

Figure 5. X-ray photoelectron spectroscopic analysis for the Pd element. The individual deconvoluted peaks of Pd 3d_{5/2} and Pd 3d_{3/2} doublet of Pd 3d spectrum are also shown. (CasaXPS, Version 2.3.22, \url{http://www.casaxps.com}).

Figure 6. (a) Frequency dependencies of dielectric constant (\varepsilon) of PbPd_{0.3}Ti_{0.7}O_3 thin films measured at different temperatures (Inset: dissipation factor (tan \delta) with frequencies at different temperatures). (b) Temperature dependent dielectric constant (\varepsilon) of PbPd_{0.3}Ti_{0.7}O_3 thin films. (Inset: dissipation factors (tan \delta) with temperatures).
polarization of $27 \mu \text{Ccm}^{-2}$, larger than those measured in other Pb-based bulk counterparts$^{19}$. The coercive field $E_c$ was obtained as $55 \text{kVcm}^{-1}$, which is also larger compared to the Pd-based ceramics ($6.5 \text{kVcm}^{-1}$)$^{19}$. In thin films, defects such as oxygen vacancies, interface layers and several other impurities generate an internal electric field$^{43,44}$. The observed discontinuity in the hysteresis loops is a result of the polarization relaxation due to this field. Such discontinuities in $P$-$E$ loops were also reported in other ferroelectric materials, e.g. $(\text{Ba}_{0.955}\text{Ca}_{0.045})_{(\text{Zr}_{0.17}\text{Ti}_{0.83})}\text{O}_3$ and Pb-based ferroelectric films$^{43,44}$. Since the structure has a polar axis, a definite proof of ferroelectricity is later confirmed from our piezoresponse force microscopy experiments.

Figure 8 shows the typical leakage current conduction behavior i.e. current-voltage ($I$-$V$) characteristic curve of the PbPdT thin films, measured with a voltage step of 0.1 V and elapsed time of 0.5 s at each voltage step. Two distinct regions are noticed in the $I$-$V$ curve: below 0.11 MV/cm, the current increases linearly with applied electric field suggesting an Ohmic conduction behavior, and above 0.11 MV/cm, the current increases exponentially with increasing electric field, which is reminiscent of Schottky or Poole-Frankel emission type conduction processes$^{45}$. A significant leakage current ($>0.001 \text{A}$) through the films is seen above 0.4 MV/cm, and it can be attributed to conductivity that originates from oxygen vacancies, and interface limited conduction processes$^{25}$. Similar conduction behavior was observed in PLD based $(\text{Ba}_{0.955}\text{Ca}_{0.045})_{(\text{Zr}_{0.17}\text{Ti}_{0.83})}\text{O}_3$ thin film capacitors, exhibiting a large leakage current density above an applied field of 0.7 MV/cm$^{45}$.

To study the ferroelectricity at nanoscale, PFM measurements were carried out on the surfaces of PbPdT films. Conducting Pt/Ir cantilever tip was used as PFM tip and it acted as the top electrode, and LSMO as the bottom electrode was grounded to obtain the piezo response images. A strong domain switching response was inferred from the amplitude and phase images (Fig. 9). A large square area of $(6 \times 6) \mu \text{m}^2$ and other central square area of $(4 \times 4) \mu \text{m}^2$ of the thin films were polled with $+12 \text{V}$ and $-12 \text{V}$ bias dc voltages, respectively. The domains with opposite polarization were distinguished with dark and bright contrast in the image. Two color contrast polling regions are clearly seen indicating the polar switching behavior in the PbPdT films in nanoscale, and the ferroelectric piezoelectric nature of the films is confirmed as that inferred from our dielectric and ferroelectric studies.
The optical absorption spectra of the thin films were measured employing a UV-2400PC spectrophotometer in the wavelength range of 300–800 nm. The substrate contribution was excluded from the spectrum by using a similar ITO-coated glass substrate as a reference. The UV-visible spectrum (Fig. 10(a)) of the PbPdT thin films show an absorption edge at 390 nm. From the Tauc plot of $(\alpha h\nu)^{2}$ versus energy (Fig. 10(b)), the direct band gap $E_g$ was calculated as 3 eV, which is found to be less than that reported for PbTiO$_3$ thin films$^{18,46}$. Thus, a reduction in band gap upon Pd substitution is apparent.

Figure 11 shows the magnetization hysteresis curves for the PbPdT films at various temperatures. The films exhibited clear magnetic hysteresis loops indicating ferromagnetism in the films. At the lowest temperature of 5 K, the saturation magnetization ($M_s$) and the coercive field ($H_c$) of the films were 6 emu·cm$^{-3}$ and 242.7 Oe, respectively. DFT calculations suggest that the substitution of non-isoelectric Pd$^{4+}$ cation at the Pb$^{2+}$ site of PbTiO$_3$ generates two electrons, which occupy the induced defect states in the electronic band gap of PbTiO$_3$. These electrons are believed to produce magnetism$^{19}$. Our XPS studies confirmed the presence of mixed oxidation states of Pd ions (Pd$^{3+}$, Pd$^{4+}$) that fulfilled the condition for magnetization in the films. To experimentally demonstrate that the Pd$^{4+}$ is at the A-site, experimental works providing local bonding coordination are required, which we are planning to carry out separately. The temperature dependence of the magnetization was studied by carrying out measurements of magnetization hysteresis loops ($M$-$H$ curve) at different temperatures (Fig. 11). A plot of the coercive field $H_c$ and the remanent magnetization $M_r$ of PbPdT films as a function of temperature is inset in Fig. 11. Upon increasing temperature, $M_r$ and $H_c$ values of the hysteresis loops decrease. These $H_c$ and $M_r$ values obtained at different temperatures are presented in Table 2. The $H_c$ value at 300 K (114 Oe) is about 2 times less than that observed at 5 K (242 Oe). Similarly, the $M_r$ value at 300 K (3.4 emucm$^{-3}$) reduces by nearly half compared to that at 5 K (6 emucm$^{-3}$). The appearance of magnetic hysteresis loop in the films at 395 K indicates that the films are ferromagnetic with a Curie temperature above 395 K. The small discontinuity in hysteresis loop

---

**Figure 9.** Piezoresponse force microscopy: (a) amplitude, and (b) phase images of PbPdT/LSMO/LSAT thin films of thickness 300 nm. The large square area of $(6 \times 6)$ $\mu$m$^2$ and central square area of $(4 \times 4)$ $\mu$m$^2$ of the thin films were polled with $+12$ V and $-12$ V bias dc voltages, respectively.

**Figure 10.** (a) UV-Visible absorption spectrum between absorbance and wavelength, (b) Optical band gap $E_g$ of PbPd$_{0.3}$Ti$_{0.7}$O$_3$ thin films using UV-Visible absorbance spectrum using Tauc plot.

---

The optical absorption spectra of the thin films were measured employing a UV-2400PC spectrophotometer in the wavelength range of 300–800 nm. The substrate contribution was excluded from the spectrum by using a similar ITO-coated glass substrate as a reference. The UV-visible spectrum (Fig. 10(a)) of the PbPdT thin films show an absorption edge at 390 nm. From the Tauc plot of $(\alpha h\nu)^{2}$ versus energy (Fig. 10(b)), the direct band gap $E_g$ was calculated as 3 eV, which is found to be less than that reported for PbTiO$_3$ thin films$^{18,46}$. Thus, a reduction in band gap upon Pd substitution is apparent.

Figure 11 shows the magnetization hysteresis curves for the PbPdT films at various temperatures. The films exhibited clear magnetic hysteresis loops indicating ferromagnetism in the films. At the lowest temperature of 5 K, the saturation magnetization ($M_s$) and the coercive field ($H_c$) of the films were 6 emu·cm$^{-3}$ and 242.7 Oe, respectively. DFT calculations suggest that the substitution of non-isoelectric Pd$^{4+}$ cation at the Pb$^{2+}$ site of PbTiO$_3$ generates two electrons, which occupy the induced defect states in the electronic band gap of PbTiO$_3$. These electrons are believed to produce magnetism$^{19}$. Our XPS studies confirmed the presence of mixed oxidation states of Pd ions (Pd$^{3+}$, Pd$^{4+}$) that fulfilled the condition for magnetization in the films. To experimentally demonstrate that the Pd$^{4+}$ is at the A-site, experimental works providing local bonding coordination are required, which we are planning to carry out separately. The temperature dependence of the magnetization was studied by carrying out measurements of magnetization hysteresis loops ($M$-$H$ curve) at different temperatures (Fig. 11). A plot of the coercive field $H_c$ and the remanent magnetization $M_r$ of PbPdT films as a function of temperature is inset in Fig. 11. Upon increasing temperature, $M_r$ and $H_c$ values of the hysteresis loops decrease. These $H_c$ and $M_r$ values obtained at different temperatures are presented in Table 2. The $H_c$ value at 300 K (114 Oe) is about 2 times less than that observed at 5 K (242 Oe). Similarly, the $M_r$ value at 300 K (3.4 emucm$^{-3}$) reduces by nearly half compared to that at 5 K (6 emucm$^{-3}$). The appearance of magnetic hysteresis loop in the films at 395 K indicates that the films are ferromagnetic with a Curie temperature above 395 K. The small discontinuity in hysteresis loop
could be due to a few missing data points in our measurements due to the lock-in amplifier (or other electronics) changing the range of measurements in the PPMS system. To examine if there are any changes in the magnetic ordering of the films, we have investigated the temperature evolution of magnetization from 5 to 395 K. A saturation field of 2 T was applied at 5 K, and the field was turned off thereafter. In the absence of the field ($H = 0$), thermal evolution of magnetization was measured up to 395 K (Fig. 12). A gradual reduction of magnetization with increasing temperature is noticed. The rate of change of magnetization with temperature ($dM/dT$) that provides information about any changes in the magnetic ordering is inset in Fig. 12. In principle, any change in magnetic ordering would result in a significant change in the magnetization curve. Since $dM/dT$ remains constant, it is evident that there is no change in magnetic ordering in the films up to 395 K. Thus, the films are ferromagnetic throughout the temperature range of the present studies and the magnetic Curie temperature is above 395 K.

Figure 11. $M$-$H$ curves measured at 5 K, 300 K and 395 K. Inset: the temperature dependence of $M_r$ and $H_c$ (lower panel).

Figure 12. Temperature evolution of the remanent magnetization created at 5 K with applied saturation field of 2 T, for PbPd$_{0.3}$Ti$_{0.7}$O$_3$ films. The inset figure shows the derivative ($dM/dT$) of magnetization $M$ with temperature $T$.

| Temperature (K) | $M_r$ (emu/cm$^3$) | $H_c$ (Oe) |
|-----------------|-------------------|------------|
| 5               | 6.007             | 242.729    |
| 10              | 5.529             | 229.906    |
| 100             | 4.052             | 165.788    |
| 200             | 3.853             | 127.318    |
| 300             | 3.455             | 114.495    |
| 360             | 3.256             | 101.671    |
| 395             | 3.162             | 96.862     |

Table 2. Remanent magnetization $M_r$ and coercive field $H_c$ at different temperatures (5–395 K) for PbPd$_{0.3}$Ti$_{0.7}$O$_3$ films.
Conclusions

(001)-oriented PbPd$_{0.6}$Ti$_{0.4}$O$_3$ thin films were grown on LSAT single crystal substrates using pulsed laser deposition technique. X-ray diffraction analysis indicates that the grown thin films are phase pure and stabilized in tetragonal phase. Temperature dependent dielectric studies on Pt/PbPdTi/LSMO metal-insulator-metal capacitors suggest that the ferroelectric Curie temperature is above 670 K. The polarization hysteresis loops at room temperature were attributed to its ferroelectric behavior. The P-E loop of our thin films showed a maximum of 92 μC/cm$^2$ polarization with remanent polarization $P_r$ of 27 μC/cm$^2$ at an applied electric field of 500 kV/cm. A saturated magnetization $M-H$ loop with remanent magnetization $M_r$ of 3.5 emu/cm$^3$ was observed at room temperature and the film retains ferromagnetic ordering up to as high as 395 K. X-ray photoelectron spectroscopic results confirmed the mixed chemical states of the Pd 3d state with Pd$^{4+}$: Pd$^{2+}$ of 0.83: 1. Magnetization in the PbPdT films is due to the existence of mixed oxidation states of Pd$^{2+}$/Pd$^{4+}$ and this fulfills the predicted requirement to realize ferromagnetism in PbPdT films. From the Tau plot of ($c0h$)2 versus energy, the direct band gap of PbPdT thin films was estimated as 3 eV. Ferroelectric piezoelectric nature of the films was confirmed from a strong domain switching response revealed from our piezoresponse force microscopy studies on the phase and amplitude contrast images. The experimental results revealed that the PbPdT thin films are multiferroic (ferroelectric-ferromagnetic) at room temperature and have potential for device applications in non-volatile memories, transducers, and actuators.

Methods

PbPdT, and PbPdT/LSMO thin films were grown on [(LaAlO)$_3$]$_n$[Sr$_2$AlTaO$_6$]$_{1-n}$ (LSAT) (001) substrates in an oxygen atmosphere using a KrF gas excimer laser operated at $\lambda = 248$ nm with a pulse repetition rate of 10 Hz using PLD technique. PbPdT target with a stoichiometric molecular formula was prepared employing the solid-state reaction method following the procedure reported earlier. The PLD deposition parameters were optimized for the growth of PbPdT thin films as summarized in Table 1. Using an XP-200 Profilometer the thickness of the films was measured to be ~300 nm. Elemental compositions and their mapping distributions were studied using Energy Dispersive X-ray Spectroscopy (EDS) employing a scanning electron microscope (JEL JSM-6480LV). Atomic Force Microscopy (AFM-Veeco) was operated in contact mode to study the surface topography of thin films and their surface roughness. The crystallographic phase and orientation of these grown thin films were confirmed from the analysis of high-resolution XRD data measured using a Smartlab x-ray diffractometer employing a Cu-Kα radiation ($\lambda = 1.5405$ Å). The chemical states of the constituent elements were studied using x-ray photo emission spectroscopy (XPS). For electrical measurements, PbPdT capacitors sandwiched between a conducting bottom LSMO layer and a top platinum (Pt) electrode were fabricated. Pt electrodes were fabricated by D.C. sputtering technique using a metal shadow mask of area $10^4$ m$^2$ followed by annealing at 620 K in an oxygen environment to recover any defects generated during sputtering. The temperature dependent dielectric properties such as capacitance and dielectric loss tangent ($\tan \delta$) were measured at several frequencies (100 Hz to 1 MHz) in the temperature range of 100–670 K by employing a programmable temperature controller (MMR K-20) with a temperature stability of $\pm 1$ K and using an impedance analyzer (HP 4294 A). Ferroelectric hysteresis loops (P-E curves) were measured at ambient temperature using a hysteresis loop tester (radiant Technologies RT6000 HVS). Magnetization measurements were carried out using a physical properties measurement system (PPMS, Quantum design), operated in the VSM module in the temperature range of 5–395 K. Piezoresponse force microscopy (PFM) studies on thin films were carried out at ambient temperature using a Multimode Nanoscope V (Veeco Instruments). The used conductive tips were coated with Pt/Ir having a resonance frequency of 140 kHz.

Data availability

The data sets are available from the corresponding author on reasonable request.

Received: 10 September 2019; Accepted: 23 December 2019; Published online: 19 February 2020

References

1. Kimura, T. et al. Magnetic control of ferroelectric polarization. Nature 426, 55 (2003).
2. Eerenstein, W., Mathur, N. D. & Scott, J. F. Multiferroic and magnetoelectric materials. Nature 442, 759 (2006).
3. Tokura, Y., Seki, S. & Nagaosa, N. Multiferroics of spin origin. Rep. Prog. Phys. 77, 076501 (2014).
4. Scott, J. F. Data storage: Multiferroic memories. Nature 442, 55 (2006).
5. Hill, N. A. Why Are There so Few Magnetic Ferroelectrics? J. Phys. Chem. B 104, 6694 (2000).
6. Wang, Y., Hu, J., Lin, Y. & Nan, C.-W. Multiferroic Magnetoelectric Composite Nanostructures. NPG Asia Mater. 2, 61 (2010).
7. Ma, J., Hu, J., Li, Z. & Nan, C.-W. Recent Progress in Multiferroic Magnetoelectric Composites: from Bulk to Thin Films. Adv. Mater. 23, 1062 (2011).
8. Vaz, C. A. F., Hoffman, J., Ahn, C. H. & Ramesh, R. Magneto-electric Coupling Effects in Multiferroic Complex Oxide Composite Structures. Adv. Mater. 22, 2908 (2010).
9. Rao, C. N. R., Sundaresan, A. & Saha, R. Multiferroic and Magnetoelectric Oxides: The Emerging Scenario. J. Phys. Chem. Lett. 3, 2237 (2012).
10. Wang, J. et al. Epitaxial BiFeO$_3$, Multiferroic Thin Film Heterostructures. Science 299, 1719 (2003).
11. Zhao, T. et al. Electrical control of antiferromagnetic domains in multiferroic BiFeO$_3$ films at room temperature. Nat. Mater. 5, 823 (2006).
12. Lebeugle, D. et al. Room-temperature coexistence of large electric polarization and magnetic order in BiFeO$_3$ single crystals. Phys. Rev. B 76, 024116 (2007).
13. Song, S. et al. Ferroelectric polarization switching with a remarkably high activation energy in orthorhombic GaFeO$_3$ thin films. NPG Asia Materials 8, e242 (2016).
14. Nechache, R. et al. Coexistence of antiferromagnetic and ferromagnetic orders at remanent state in epitaxial multiferroic Bi₅FeCrO₉ nanostructures. J. Phys.: Condens. Matter. 24, 222002 (2012).
15. Quattropani, A. et al. Band–Gap Tuning in Ferroelectric Bi₅FeCrO₉ Double Perovskite Thin Films. J. Phys. Chem. C 122, 1070 (2018).
16. Nechache, R. et al. Epitaxial thin films of multiferroic Bi₅FeCrO₉ with B-site cationic order. J. Mater. Res. 22, 2102 (2007).
17. Catalan, G. & Scott, J. F. Physics and Applications of Bismuth Ferrite. Adv. Mater. 21, 2463 (2009).
18. Mansoor, M. A. et al. Perovskite-Structured PbTiO₃ Thin Films Grown from a Single-Source Precursor. Inorg. Chem. 52, 5624 (2013).
19. Kumari, S. et al. Palladium-based ferroelectrics and multiferroics: Theory and experiment. Phys. Rev. B 95, 214109 (2017).
20. Guthrie, A. N. & Copley, M. J. The Magnetic Moment of the Palladium Atom. J. Appl. Phys. Status Solidi B 124, 1403 (2014).
21. Obinata, A. et al. Electric-field control of magnetic moment in Pd. Scientific Reports 5, 14303 (2015).
22. Bursan, M., Fulz, B. & Anthony, L. Magnetic properties of bcc Fe–Pd extended solid solutions. Phys. Rev. B 55, 11502 (1997).
23. Evans, D. M. et al. The Nature of Magnetoelastic Coupling in Pb(Zr,Ti)O₃. Adv. Mater. 27, 6068 (2015).
24. Bennett, J. W., Grinberg, I., Davies, Peter, K. & Rappe, A. M. Pb-free semiconductor ferroelectrics: A theoretical study of Pd-substituted Ba(Ti, Ce)O₃ solid solutions. Phys. Rev. B 82, 184106 (2010).
25. Graduautea, E. et al. Lead palladium titanate: A room-temperature multiferroic. Phys. Rev. B 96, 104104 (2017).
26. Eerenstein, W. et al. Comment on Epitaxial BiFeO₃ Multiferroic Thin Film Heterostructures. Science 307, 1203 (2005).
27. Levstik, A. et al. Magnetoelastic relaxor. Appl. Phys. Lett. 91, 012905 (2007).
28. Kumar, A., Rivera, I., Katiyar, R. S. & Scott, J. F. Multiferroic Pb(Fex0.75W0.25)TiO₃ thin films: A room-temperature relaxor ferroelectric and weak ferromagnetic. Appl. Phys. Lett. 92, 132913 (2008).
29. Ortega, N., Kumar, A., Bhattacharya, P., Majumder, S. B. & Katiyar, R. S. Impedance spectroscopy of multiferroic PbZr(Ta, Ti, O3–x)O/CoFe2O4 layered thin films. Phys. Lett. B 77, 014111 (2008).
30. Schrom, D. G., Chen Long-Qing, P., Xiaqing Schmehl, A. & Zurzurbuch, M. A. J. Am. Ceram. Soc. 91, 2429–2454 (2008).
31. Chrisey, D. B. & Hubler, G. K. Pulsed Laser Deposition of Thin Films. 1st edn (Wiley-Interscience, 1994).
32. Bhattarai, M. K., Mishra, K. K., Dugu, S., Instan, A. A. & Katiyar, R. S. Ferroelectric ordering and energy storage density of thin films capacitor by doping La⁺[2] and Sc⁺[2] on Pb(Zr0.65,Ti0.35)O₃ using pulsed laser deposition technique. Appl. Phys. Lett. 114, 223902 (2019).
33. Bhattarai, M. K., Mishra, K. K., Instana, A. A., Bastakoti, B. P. & Katiyar, R. S. Enhanced energy storage density in Sc⁺[2] substituted Pb(Zr0.65,Ti0.35)O₃ nanoscale films by pulse laser deposition technique. Appl. Surf. Sci. 490, 451 (2019).
34. Barrionuevo, D. et al. Thickness dependent functional properties of PbZr0.65Ti0.35O3/La0.65Sr0.35MnO3 heterostructures. J. Appl. Phys. 114, 234103 (2013).
35. Majumdar, S., Huhtinen, H., Paturi, P. & Majumdar H. S. Effect of strain and grain boundaries on dielectric properties in La₅Sr₅MnO₁₅ thin films. J. Mater. Sci. 48, 2115 (2013).
36. Kumar, S. et al. Effect of thickness on dielectric, ferroelectric, and optical properties of Ni substituted Pb(Zr₀.₅₅Ti₀.₄₅)O₃ thin films. J. Appl. Phys. 118, 184103 (2015).
37. Bajpai, P. K., Pastor, M. & Singh, K. N. Diffuse Phase Transition and Electrical Conductivity of Pb(Ca₁₋ₓNbₓ)O₃. J. Electron. Mater. 43, 1403 (2014).
38. Singh, G., Tiwari, V. S. & Gupta, P. K. Role of oxygen vacancies on relaxation and conduction behavior of KNbO₃ ceramic. J. Appl. Phys. 107, 064103 (2010).
39. Newham, R. E. Properties of Materials: Anisotropy, Symmetry. Structure. Kindle edn (Oxford University Press Inc., 2005).
40. Ryu, J. et al. Upshift of Phase Transition Temperature in Nanostructured PbTiO₃ Thick Film for High Temperature Applications. ACS Appl. Mater. Interfaces 6, 11990 (2014).
41. Singh, U. G. et al. Pd-doped perovskite catalyst, BaCe₂₋₀.₅Pd₀₋₀.₅O₉₋₋, for CO oxidation. J. Catal. 249, 347 (2009).
42. Zhu, Y. et al. Boosting Oxygen Reduction Reaction Activity of Palladium by Stabilizing Its Unusual Oxidation States in Perovskite. Chem. Mater. 27, 3048 (2015).
43. Thomas, R., Mochizuki, S., Mihara, T. & Ishida, T. Preparation of Pb(Zr,Ti)O₃ thin films by RF–magnetron sputtering with single stoichiometric target: structural and electrical properties. Thin Solid Films 413, 65 (2002).
44. Lee, J. & Ramesh, R. Imprint of (PbLa)(Zr,Ti)O₃ thin films with various crystalline qualities. Appl. Phys. Lett. 68, 484 (1996).
45. Puli, V. S. et al. Nanoscale polarisation switching and leakage currents in (Ba₀.₈₅Ca₀.₁₅)(Zr₀.₇₅Ti₀.₂₅)O₃ epitaxial thin films. J. Phys. D: Appl. Phys. 48, 355502 (2015).
46. Zanetin, V. I. Absorption Edge Anomalies in Polar Semiconductors and Dielectrics at Phase Transitions. Phys. Status Solidi B 124, 625 (1984).
47. Momma, K. & Izumi, F. VESTA 3 for three-dimensional visualization of crystal, volumetric and morphology data. J. Appl. Crystallogr. 44, 1272 (2011).

Acknowledgements
The authors acknowledge the financial support from the Department of Defense, USA (DoD Grant No. FA9550-16-1-0295). K.K.M thanks to all SPECLAB members for their help in experimental studies.

Author contributions
K.K.M synthesized and characterized the thin films. The data were analyzed by K.K.M. A.A.I. involved in the PFM analysis. S.K. helped with the XPS analysis. R.S.K and J.F.S supervised this research work. Manuscript written by all co-authors co-wrote and commented on the manuscript.

Competing interests
The authors declare no competing interests.

Additional information
Correspondence and requests for materials should be addressed to K.K.M. or R.S.K.

Reprints and permissions information is available at www.nature.com/reprints.

Publisher’s note
Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.
