Magnetic properties of Gd\(T_2\)Zn\(20\) (\(T = \text{Fe, Co}\)) investigated by X-ray diffraction and spectroscopy

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We investigate the magnetic and electronic properties of the Gd\(T_2\)Zn\(20\) (\(T = \text{Fe and Co}\)) compounds using the X-ray resonant magnetic scattering (XRMS), X-ray absorption near-edge structure (XANES) and X-ray magnetic circular dichroism (XMCD) techniques. The XRMS measurements reveal that the Gd\(Co_2\)Zn\(20\) compound has a commensurate antiferromagnetic spin structure with a magnetic propagation vector \(\vec{q} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})\) below the Néel temperature (\(T_N \sim 5.7\) K). Only the Gd ions carry a magnetic moment forming an antiferromagnetic structure with magnetic representation \(\Gamma_6\). For the ferromagnetic GdFe\(2\)Zn\(20\) compound, an extensive investigation was performed at low temperature and under magnetic field using XANES and XMCD techniques. A strong XMCD signal of about 12.5 % and 9.7 % is observed below the Curie temperature (\(T_C \sim 85\) K) at the Gd-L\(2\) and L\(3\) edges, respectively. In addition, a small magnetic signal of about 0.06 % of the jump is recorded at the Zn K-edge suggesting that the Zn 4\(p\) states are spin polarized by the Gd 5\(d\) extended orbitals.

I. INTRODUCTION

The demand for new materials with interesting and useful physical properties has led to a fast development in material science. Properties such as superconductivity, magnetic ordering, nearly ferromagnetic Fermi-liquid and heavy fermion behavior have been observed in many materials, including the family of complex intermetallic compounds \(RT_2\)Zn\(20\) (\(R = \text{rare earth, } T = \text{transition metal}\)). This family, first reported two decades ago by Nasch et al.6 has been extensively used as a model system due to its rather unique structure which features a complex but well ordered crystal structure. The \(RT_2\)Zn\(20\) compounds have a cubic structure with \(Fd\bar{3}m\) (No. 227) space group in which the \(R\) and \(T\) ions occupy the 

FIG. 1. (Color online) \(RT_2\)Zn\(20\) structural representation. The Zn ions are shown at the three sites (label Zn1-3) and the Zn cages are shown in detail in which there is a \(T\) atom inside.

transition temperature. When the \(Co\) ions are replaced by Fe ions (GdFe\(2\)Zn\(20\) and TbFe\(2\)Zn\(20\)) the compounds exhibit ferromagnetic (FM) ordering with the transition temperatures drastically raised to 86 K and 66 K, respectively. The relatively long distance between rare earth ions in the structure weakens the Ruderman-Kittel-Kasuya-Yosida (RKKY) exchange interaction, while the high density of Fe 3\(d\) bands at the Fermi level directly affects the conduction electrons.

Recently, a detailed study of the magnetic structure of TbCo\(2\)Zn\(20\) and TbFe\(2\)Zn\(20\) using magnetic neutron scattering at low temperature was reported.2 Although the Tb-based compounds present similar magnetic properties as compared to GdCo\(2\)Zn\(20\) and GdFe\(2\)Zn\(20\), the absence of crystalline electric field (CEF) at first order and a strong RKKY interaction in the Gd-compounds affect the electronic and magnetic interactions between the rare earth ions and the surrounding matrix. In particular, the investigation of the compounds with half-filled 4\(f\)\(^2\) shell (Gd-based materials) at low temperature can...
provide information about the physical properties and are very important as reference compounds due to their lack of orbital momentum, i.e., $S = 7/2$ and $L = 0$, which leads to magnetic properties that are unaffected by spin-orbit coupling. In order to understand the implications of such interactions we have used spectroscopy and magnetic scattering techniques to probe in detail the electronic and magnetic properties of the Gd$T_2$Zn$_{20}$ family. Due to the large Gd neutron absorption cross section, X-ray technique is the ideal alternative to probe the magnetic and electronic properties in these materials. Furthermore, the incoming beam energy can be tuned to the absorption edge, thus providing chemical and atomic selectivity, i.e., the magnetic response of each element can be probed separately.

Here we report the magnetic and electronic structure at low temperature of the Gd$T_2$Zn$_{20}$ ($T = $ Fe$_{20}$ Co) compounds using the X-ray resonant magnetic scattering (XRMS), X-ray absorption near-edge structure (XANES) and X-ray magnetic circular dichroism (XMCD) techniques. The XRMS measurements performed on GdCo$_2$Zn$_{20}$ reveal a commensurate antiferromagnetic ordering with a magnetic propagation vector $\mathbf{q} = (1/2, 1/2, 1/2)$ below $T_N = 5.72(6)$ K. In addition, at low temperature the Gd magnetic moments order following the magnetic representation $\Gamma_6$ in which the magnetic moment direction is written as a linear combination of the two basis vectors $\psi_3$ and $\psi_4$. This magnetic structure is consistent with a $P_3\bar{1}$ magnetic space group. The XMCD measurements performed below the Curie temperature ($T_C = 85(2)$ K) in GdFe$_2$Zn$_{20}$ display a dichroic signal of 12.5 % and 9.7 % of the absorption jump for Gd $L_2$ and $L_3$ edges, respectively. Surprisingly, a magnetic signal of about 0.06 % is detected at the Zn $K$-edge which suggests that the Zn ions are spin polarized. This magnetic signal might originate from the hybridization between the extended Gd 5$d$ bands with the empty Zn 4$s$ states. Absorption measurements performed at the Fe $K$-edge do not reveal any magnetic contribution coming from the iron ions above the background level.

II. EXPERIMENTAL DETAILS

High quality single crystals of GdFe$_2$Zn$_{20}$ and GdCo$_2$Zn$_{20}$ were grown at UFABC by Zn self flux method, similar to that reported in previous studies on the family. In order to perform the absorption measurements at the Gd $L_{2,3}$, Fe and Zn $K$ edge, selected single crystals of GdFe$_2$Zn$_{20}$ were ground and sieved, resulting in fine powders with grain sizes around 3-5 $\mu$m. The magnetic diffraction measurements were done on 4-ID-D at the Advanced Photon Source (APS), Argonne National Laboratory (Argonne, IL/USA), whereas the absorption measurements were conducted at 4-ID-D (APS) and at beamline P09 at PETRA III (DESY, Hamburg/Germany).

A. Absorption experiment

XANES and XMCD spectra obtained at low temperature for the Gd $L_{2,3}$, Fe and Zn $K$-absorption edges were performed in transmission geometry on powdered GdFe$_2$Zn$_{20}$ samples. The samples were cooled down by a diplex cryostat with base temperature around 7 K. XMCD spectra were performed in helicity switching mode in which the left and right circular polarization was obtained by means of diamond phase plates. The degree of circularly polarized beam was higher than 95 % for both beamlines (P09 and 4-ID-D). An external magnetic field of $H = 2.0$ T (at APS) and 0.8 T (at DESY) was applied in the GdFe$_2$Zn$_{20}$ samples along and opposite to the incident beam wave vector $k$ to align the ferromagnetic domains and to correct for non-magnetic artifacts in the XMCD data. Those external magnetic fields were enough to reach the saturation magnetization according to the macroscopic measurements.

B. Scattering experiment

XRMS measurements were performed at $T = 4.5$ K on GdCo$_2$Zn$_{20}$ single crystal, mounted inside the closed-cycle diplex cryostat in a six-circle diffractometer at the 6-ID-B beamline. The single crystal was oriented with the [111] direction parallel to the vertical diffraction plane. Several magnetic superlattice reflections of the type $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ with $L = (2n+1)$ were measured and their integrated intensities were compared to the simulated intensities to determine the magnetic structure below $T_N$. To enhance the magnetic Bragg peak intensities, the energy of the incident beam was tuned near the Gd $L_3$ or $L_3$ absorption edges. In addition, in order to investigate any magnetic contribution from Co and Zn ions, the energy of the incident beam was also tuned to the Co (7709 eV) and Zn (9659 eV) $K$-edges and a search for superlattice reflections was performed. The charge and magnetic contributions present in the scattered beam were separated by a pyrolytic graphite [C(006)] analyzer crystal installed on the 20 arm of the diffractometer. Since the incident beam presents the polarization perpendicular to diffraction plane ($\sigma$ polarization), by rotation of the analyzer...
crystal around the scattered beam wave vector $k'$ we were able to select the two polarization channels ($\sigma - \sigma'$ and $\sigma - \pi'$) in this experimental geometry.\textsuperscript{17}

III. EXPERIMENTAL RESULTS

The experimental results are organized into two sections: The first part is dedicated to the absorption measurements on powdered GdFe\textsubscript{2}Zn\textsubscript{20} in its FM state. The subsequent section shows the results obtained by the XRMS technique on the GdCo\textsubscript{2}Zn\textsubscript{20} single crystal in its AFM state.

A. Absorption measurements: GdFe\textsubscript{2}Zn\textsubscript{20}

XANES and XMCD measurements performed at the Gd $L_{2,3}$ edges in GdFe\textsubscript{2}Zn\textsubscript{20} are shown in Fig. 2. Dipolar selection rules make the dichroic signal at the $L_{2,3}$ absorption edges sensitive to the spin polarization of the $\gamma$ intermediate 5$d$ level. The Gd XANES reported in Fig. 2a are normalized to one at the $L_3$ and half at the $L_2$ edge to reflect the 2:1 ratio of the initial state at these edges\textsuperscript{40} ($2p_{3/2}$ and $2p_{1/2}$, respectively). Figure 2 also shows the XMCD spectra at the Gd $L_{2,3}$ edges in which each spectrum is normalized to the corresponding edge jump of the absorption spectrum. The XMCD signal obtained across the two edges show different intensities with a strong dichroic magnetic signal around 12.5% at the $L_2$ and 9.7% at the $L_3$ absorption edges, which is consistent with Gd-based compounds.\textsuperscript{25} The size and the shape of the magnetic contribution obtained by fitting the XMCD signals with Lorentzian function can describe additional properties of this system. The widths of the dipolar contributions (E1) contributions observed at the $L_2$ and $L_3$ absorption edges are 4.3(2) eV and 4.5(2) eV, respectively, which reflects a short 2$p$ core hole lifetime. Using the integrated intensities, the $L_3/L_2$ ratio (or Branching Ratio value - BR)\textsuperscript{19-23} obtained experimentally is 0.77(4).

Absorption measurements at the Fe and Zn $K$ edges were also carried out on powdered samples. The absorption measurement near the $K$ edge, in which the dipolar transition is probed ($1s \rightarrow 4p$), is crucial towards understanding the magnetic and electronic properties due to the delocalized character of the $p$ states.\textsuperscript{24} Since the probed $p$ states are very delocalized, a strong influence of the surrounding matrix can be expected due to the hybridization between the rare earth and the transition metal ions. As shown in Figure 3, the measurements performed near the Fe $K$ edge do not reveal any magnetic contribution from the Fe ions higher than the background level ($\sim 0.07\%$). The inset in Fig. 3 exhibits the XMCD measurements obtained for a 5 $\mu$m Fe-foil in the same experimental conditions: a clear dichroic signal can be observed near the edge.

However, the spectroscopy measurement at the Zn $K$ edge manifests an interesting behavior. Figure 4 shows the absorption and dichroism results at the Zn $K$ edge in which an induced magnetic signal around 0.06% is detected. This magnetic signal is due to a hybridization...

FIG. 2. (Color online) XANES and XMCD spectra obtained at the Gd $L_{2,3}$ edges performed at $T = 7$ K and $H = 2$ T.

FIG. 3. (Color online) X-ray absorption measurements at 7 K and in 2 T at the Fe $K$-absorption edge for the GdFe\textsubscript{2}Zn\textsubscript{20}. The inset shows the XMCD data for the Fe-foil sample.

FIG. 4. (Color online) X-ray absorption measurements at 7 K and in 2 T for Zn $K$-absorption edge.
with the rare earth 5d orbitals. The XMCD spectrum exhibits the main positive feature located at 9665 eV with a width around 2.5 eV, surrounded by two negative peaks 6.5 eV away. The broad feature localized around 9680 eV (∼20 eV above the edge) is likely due to magnetic EXAFS.

A clear evidence of the Zn 4p states polarization due to the Gd ions can be found in the temperature and field dependence reported in panels (a) and (b), respectively. The magnetic intensities for both Gd and Zn ions follow the same temperature evolution and disappear around the critical temperature (T_C ∼ 85 K). In addition, the two hysteresis loops obtained at the maximum XMCD intensity shows clearly the Zn magnetism dependence in relation to the Gd ions and therefore it suggests a spin polarization of the Zn 4p bands by the Gd sub-lattice.

B. Magnetic scattering measurements: GdCo_2Zn

Figure 5 shows the evolution of the integrated intensity for the magnetic Bragg reflection (2 4 3 5 6) as a function of temperature for the GdCo_2Zn compound fitted by a Lorentzian-Squared function. The magnetic peak intensity decreases smoothly to zero as the temperature approaches T_N, indicating a standard second order phase transition from an AFM to a paramagnetic state. A dashed red line in Fig. 5 displays a fitting using a critical power-law expression, (1 − T/T_N)^2, above 5.0 K. The fitting around the Néel temperature yields a T_N = 5.72(6) K and a critical exponent β = 0.36(3). The value of T_N is in good agreement with bulk magnetic susceptibility measurements and previous works. The critical exponent β close to 0.367 suggests a three-dimensional (3D) Heisenberg magnetic model. In blue symbol (Fig. 6) is also reported the full width at half maximum (FWHM) of the magnetic superlattice peak (2 4 3 5 6) as a function of temperature. This figure clearly shows a peak broadening and a decrease in intensity near the phase transition temperature characteristic of a loss of long-range order. The width of the magnetic Bragg peak is inversely proportional to the correlation length (ξ). From the FWHM data, the estimated correlation length at low temperature is ξ ~ 1100 Å.

Energy dependences across the Gd-L_2 and L_3 edges performed at 4.5 K are displayed in Fig. 6. The top panels [Fig. 7(a) and 7(c)] show the normalized absorption coefficients (µ) obtained from the fluorescence yield, while the bottom panels [Fig. 7(b) and 7(d)] exhibit the energy profile obtained at the magnetic superlattice position (2 4 3 5 6) with the analyzer crystal set to the σ − π' polarization channel. A resonant enhancement of over two orders of magnitude at both absorption edges can be seen. In addition, the maximum intensities are observed about 2-3 eV above the absorption edge (defined by the vertical dashed lines), which is a characteristic signature of a dipole electronic transition. The same energy dependence was performed in the σ − π' polarization channel, and no significant contribution was observed. The strong resonant enhancement in the spectra [Fig. 7(b) and 7(d)] indicates a significant overlap between the initial 2p and 5d states, and a strong exchange interaction between the 4f-5d orbitals. This magnetic polarization of the 5d bands via 4f states helps shed light on the magnetic structures of these rare earth based materials using the L absorption edge measurements, i.e., 2p → 5d transitions. Moreover, the asymmetric peak shape expressed as a long tail below the absorption edges arises from the interference between the resonant and non-resonant magnetic scattering contributions.

The normalized energy line shape dependence after absorption correction for selected magnetic Bragg peaks (2 4 3 5 6) with L = (2n + 1) performed at Gd-L_2 and L_3.
rescence emission. The lower pictures [(b) and (d)] exhibit the L\textsubscript{Gd} edges are displayed in Fig. 7 for the space group Fd\(\bar{3}\)m:2 with \(k_0 = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})\). The decomposition of the magnetic representation (MR) for the Gd site can be written like \(\Gamma_{MAG} = \Gamma_2^L + \Gamma_3^L + \Gamma_5^L + \Gamma_8^L\). The two rare earth atoms positions of the nonprimitive basis are defined according to 1: (0.125, 0.125, 0.125) and 2: (0.875, 0.875, 0.875).

### Table I. Basis vectors (BV’s) for the group Fd\(\bar{3}\)m:2

| IR BV Atom | \(m_{||a}\) | \(m_{||b}\) | \(m_{||c}\) |
|------------|-----------|-----------|-----------|
| \(\Gamma_2\) | 1         | 1         | 1         |
| \(\psi_1\)  | 0         | 0         | 0         |
| \(\Gamma_3\) | 2         | 1         | 1         |
| \(\psi_2\)  | 0         | 0         | 0         |
| \(\Gamma_5\) | 2         | -1        | -1        |
| \(\psi_3\)  | 0         | 0         | 0         |
| \(\psi_4\)  | 1         | 0         | 0         |
| \(\psi_5\)  | 0         | 0         | 0         |
| \(\psi_6\)  | 0         | 0         | 0         |
| \(\psi_7\)  | 0         | 0         | 0         |

The magnetic structure of GdCo\(_2\)Zn\(_{20}\)

The magnetic structure of the Gd spins is determined comparing the experimental integrated intensities in Fig. 8 with simulated data from selected magnetic reflections. The SARA\textsuperscript{11} and ISODISTORT\textsuperscript{12} softwares were used to determine the possible magnetic arrangements that the Gd ions can adopt inside the unit cell, i.e., the magnetic representation (\(\Gamma_{MAG}\)) and its magnetic space group. In addition, we assumed that only the Gd ions carry magnetic moments in this compound. For this material, whose magnetic propagation vector is \((\frac{1}{2}, \frac{1}{2}, \frac{1}{2})\) (point \(L\) in the Brillouin zone), whose space group is Fd\(\bar{3}\)m, and whose Gd ions occupy the \(8a\) crystallographic site, the magnetic representation (MR) can be decomposed into four non-zero irreducible representations (irreps): two one-dimensional (1D - \(\Gamma_{1,2,3}\)) and two two-dimensional (2D - \(\Gamma_{5,6}\)). The four possible magnetic representations for the AFM GdCo\(_2\)Zn\(_{20}\) compound are summarized in Table 1 with their respective basis vectors (BV’s). The propagation vector and the irreps are labeled following the Kovalev notation\textsuperscript{33}, as given by the program SARA\textsubscript{h}.

To determine the magnetic structure, the intensities were calculated assuming only dipole transition (E1) and hence, the X-ray magnetic scattering cross section model can be written as \(17,34–38\).
The resonant term. It carries information about the magnetic configuration.

The term $A$ [Eq. (2)] contains the absorption correction and the Lorentz factor. $\alpha$ is the angle between the wave-vector transfer $\vec{Q} (= \vec{k'} - \vec{k})$ and the [111] crystal direction, and $\theta$ is half of the $2\theta$ scattering angle. Eq. (3) shows the resonant term. It carries information about the $\vec{k}$ ($\hat{\varepsilon}$) and $\vec{k'}$ ($\hat{\varepsilon}'$), i.e., the incident and scattered wave (polarization) vectors, respectively, and the magnetic moment direction at nth site ($\hat{z}_n$). The terms $F(0,1,2)$ are related to the dipole matrix transition and by atomic properties. The exponential function in Eq. (4) is a function of the wave-vector transfer $\vec{Q}$, and the position $\vec{R}_n$ of the nth Gd ion inside the unit cell. For the XRS technique probing AFM materials, the magnetic intensity at the first harmonic satellites comes from the linear term on magnetic moment direction: $[-i(\hat{\varepsilon}' \times \hat{\varepsilon}) \cdot \hat{z}_n]$ displayed in Eq. (3).

The simulated intensities obtained using Eq. (1) and the experimental intensities obtained at the Gd L$_2$ edge [Fig. 7(g-l)] are displayed in Fig. 9. The two irreps ($\Gamma_2$) are displayed in Fig. 9. The two irreps imply the magnetic spin moments aligned along the [1,1,1] crystallographic direction. On the other hand, for the $\Gamma_5$ and $\Gamma_6$ irreps, the magnetic moment can be written as a linear combination of the BV's $\psi_3$ and $\psi_4$ for $\Gamma_5$, and $\psi_5$ and $\psi_6$ for $\Gamma_6$, i.e., $\mathbf{z}_n = c_{3,n} \psi_3 + c_{4,n} \psi_4$ or $\mathbf{z}_n = c_{5,n} \psi_5 + c_{6,n} \psi_6$, with $c_{3-6}$ real or complex numbers. As displayed in Fig. 9, the simulated intensities regarding the Gd spins moment directions were performed for different magnetic representation. The better agreement (smaller $\chi^2$ of 1.5) is achieved when the magnetic structure is defined by the representation $\Gamma_6$ with $c_5 = -0.24$ and $c_6 = 0.27$. Using this information and the ISODISTORT software, we could assign the $P_{21}$ magnetic space group to this magnetic configuration.

To identify different magnetic propagation vectors, such as $(0, 0, 1/2)$ and $(0, 0, 1)$, a systematic search for commensurate and incommensurate magnetic reflections in the reciprocal space was performed below $T_N$. However, only magnetic Bragg reflections of the type $(1/2, 1/2, 1)$ were observed. In addition, to probe a possible presence of AFM magnetic moments in the Co and Zn ions, the beam energy was tuned to the Co and Zn K edges and a search for magnetic superlattice reflections was performed below $T_N$. No measurable magnetic reflections were observed at 4.5 K could be observed at those edges.

**IV. DISCUSSION**

The nature of the electronic and magnetic properties of the intermetallic RT$_2$Zn$_20$ systems depends strongly on the interactions between the rare earth and the transition metal ions. Since the Gd 5d states participate on the conduction bands, the resonant absorption and diffraction measurements at the Gd L$_2,3$ edges provide valuable information. As reported in susceptibility measurements and band structure calculations, the replacement of the transition metal affects the electronic density of states at the Fermi level ($\rho_{E_F}$) and the conduction band without changing significantly the lattice parameters. Doping studies of Gd(Fe$_x$Co$_{1-x}$)$_2$Zn$_{20}$ compounds shows a monotonic increase of the magnetic ordering temperature when $x$ increases to 1 which indicates that there is an increase in the coupling between the rare earth magnetic moments.

Absorption measurements performed in powdered samples of FM GdFe$_2$Zn$_{20}$ compound below $T_C$ reveal interesting behaviors. As can be seen in Fig. 3, only at the Gd $L$ and Zn $K$-edges a dichroic signal is observed above the background level whilst an unexpected lack of magnetic intensity is observed at the Fe $K$-edge. The intense magnetic signal at the rare earth $L$-edges occurs mainly due to the overlap between the Gd 2$p$ and 5$d$ states and a strong energy splitting of the 5$d$ subbands as a result of a 4$f$-$5d$ exchange interaction. In addition, the splitting of the $d$ states into 5$d$ spin-up and spin-down wave functions has considerable influence on the magnetism observed at the Zn $K$-edge. The Zn 3$d$ orbitals are completely filled (3$d^{10}$) and henceforth, a magnetic moment due to an overlap between the 3$d$
and 4sp orbitals in the Zn ion is not expected to occur due to the filled 3d orbitals being more contracted. Therefore, the magnetic signal observed in the 4p states is due to hybridization with the extended Gd 5d orbitals and not from the exchange interaction with the Zn 3d\textsuperscript{10} orbitals. Following Hund’s rule, the Gd ion has the 4f state filled by seven spin-up electrons in which it pulls the 5d\textsubscript{sub band} spin-up function towards the inner core due to a positive exchange interaction. The short distances between the first Gd-Zn ions (≈ 3 Å) drives a quenching of the angular momentum, so the 5d\textsubscript{3/2} and 5d\textsubscript{5/2} sub-bands should display the same polarization and thereby the dichroism at the L\textsubscript{3} and L\textsubscript{2} edges should have equal magnetic intensity. For the GdCo\textsubscript{2}Zn\textsubscript{20} compound, the BR ratio varies between 0.9 - 1.2, i.e., close to the theoretical value and therefore we cannot suggest that this magnetic moment must be reported by band structure calculation. The magnetic moment of the Fe ions is also observed in the TbCo\textsubscript{2}Zn\textsubscript{20} materials, the crystalline electric field (CEF) splitting can affect the magnetic representation \(\Gamma_5\) in which the magnetic moment directions were performed for different magnetic field. The absence of magnetism in the transition metal ions affects the environment around the atoms and consequently spin polarizes the transition metal ion. The rough estimate of magnetic moment for the Zn ions based on spectroscopy and macroscopic measurements has to be further investigated.

To provide further information about these systems, the magnetic properties of the AFM compound GdCo\textsubscript{2}Zn\textsubscript{20} were also investigated, using the XRMS technique. The transition to the magnetically ordered phase driven by temperature is characterized by the appearance of superlattice magnetic reflections with a magnetic propagation vector \(\vec{q} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})\). This \(\vec{q}\) magnetic vector indicates that the magnetic unit cell is represented by a doubled chemical unit cell in all three crystallographic directions. In addition, as seen in Fig. 9, the experimental and simulated intensities regarding the Gd ions moment directions were performed for different magnetic representation and the magnetic structure is identified as magnetic representation \(\Gamma_5\) in which the magnetic moment can be written as a linear combination of the BV’s \(\psi_5\) and \(\psi_6\). For this magnetic structure, the magnetic space group is \(P\bar{S}1\).

Since the closest distance between the rare earth ions is \(≈ 6\) Å, the magnetic properties in this system will be mainly mediated via the conduction electrons. This large distance explains quite well a weakening of the \(J_{RKKY}\) exchange interaction and thus a very low representative magnetic transition temperature, i.e., approximately 5.7 K (GdCo\textsubscript{2}Zn\textsubscript{20}) and 2.5 K (TbCo\textsubscript{2}Zn\textsubscript{20}). As a consequence of a poor \(J_{RKKY}\) coupling, the surrounding matrix around the \(R\) ions is weakly affected and hence it is not possible to induce a magnetic moment in the transition metal ions. The absence of magnetism in the transition metal ions is also observed in the TbCo\textsubscript{2}Zn\textsubscript{20} compound. For the Tb-based materials, the crystalline electric field (CEF) splitting can affect the \(J_{RKKY}\) constant and therefore influence the magnetic coupling and the transition temperature. It is well known that the CEF splitting induces magnetic anisotropies in the ground state and that it may influence the total angular momentum. Several
Tb-based intermetallics have exhibited distinct magnetic properties in relation to the Gd counterparts as shown in the layered family $R_n T_m M_{3n+2m}$ ($R =$ rare earth; $T =$ transition metal; $n = 1,2$ and $m = 0,1,2$). Although the GdCo$_2$Zn$_{20}$ and TbCo$_2$Zn$_{20}$ compounds display the same magnetic propagation vector ($\frac{1}{2}, \frac{1}{2}, \frac{1}{2}$), their magnetic structures are different, which is mainly related to the competition between the RKKY and CEF interactions. Jia et al. evaluated the CEF parameters from the thermodynamic measurements for the entire $R$Co$_2$Zn$_{20}$ series ($R =$ Tb-Tm) and they observe small energy scales and a large $B_0^2$ CEF parameter for the complete series. This finding suggests a small energy level splitting and a strong influence of the Zn cage on the rare earth ions, i.e., guest-framework interaction. Therefore, we suggest that the rare earth ions located in this large polarized environment are strongly affected by the Zn cages, which has a direct influence on the electronic and magnetic properties. This can be seen in the different magnetic structures for the AFM compounds and the spin polarization of the Zn ions only for the GdFe$_2$Zn$_{20}$ compound. Hence, the CEF effect has an important role in this class of compound. A detailed investigation for different rare earth elements would allow a better understanding of the $R$Co$_2$Zn$_{20}$ family. Nevertheless macroscopic measurements down to 1.8 K report a magnetic ordering only for the compounds with $R =$ Gd and Tb.

V. SUMMARY

We have investigated the intermetallic Gd$T_2$Zn$_{20}$ system with $T =$ Co and Fe at low temperature using the XRS measurements and XANES/XMCD techniques. The XRMS measurements performed in GdCo$_2$Zn$_{20}$ compound reveal a commensurate antiferromagnetic ordering with a magnetic propagation vector ($\frac{1}{2}, \frac{1}{2}, \frac{1}{2}$) in which only the Gd ions carry magnetic moments. Selected magnetic reflections were measured in the polarization channel $\sigma - \pi'$ and we identified that the Gd spins follow the magnetic representation $\Gamma_6$, which is different from the isostructural compound TbCo$_2$Zn$_{20}$, mainly due to the CEF effects in the latter. The evolution of magnetic signal showed a magnetic phase transition below $T_N = 5.72(6)$ K with a critical exponent $\beta = 0.36(3)$, suggesting a three-dimensional (3D) Heisenberg magnetic model. The XANES and XMCD measurements performed at the Gd $L_2,3$ edges in GdFe$_2$Zn$_{20}$ reveal a strong magnetic signal ($\sim 12.5 \% - L_2$ and $9.7 \% - L_3$) indicating a splitting of the 5$d$ orbitals and a strong Gd-Gd exchange interaction as well as a non zero orbital moment. In addition, we observe the presence of a small magnetic dichroic signal at the Zn $K$ edge due to the spin polarization of the Gd 5$d$ orbitals. This indicates a large RKKY exchange interaction between the Gd-Gd ions which polarizes the surrounding matrix.

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The minus sign arising from the spin-orbit interaction in the initial $2p^x$ state.

The GdFe$_2$Zn$_{20}$ compound has a cubic lattice parameter $a$ around 0.07 Å larger than GdCo$_2$Zn$_{20}$. Hence, a modification in the unit cell volume $\sim 3\%$ by replacing the ion.