Global existence result for phase transformations with heat transfer in shape memory alloys*

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Abstract. We consider three-dimensional models for rate-independent processes describing materials undergoing phase transformations with heat transfer. The problem is formulated within the framework of generalized standard solids by the coupling of the momentum equilibrium equation and the flow rule with the heat transfer equation. Under appropriate regularity assumptions on the initial data, we prove the existence a global solution for this thermodynamically consistent system, by using a fixed-point argument combined with global energy estimates.

Key words. Existence result, generalized standard materials, heat equation, enthalpy transformation, maximal monotone operators, doubly nonlinear equations, shape-memory alloys.

AMS 2000 Subject Classification: 35K55, 74C05, 74C10, 74F05, 74N30, 80A17.

1 Introduction

Motivated by the study of shape memory alloys, we consider rate-independent processes describing materials undergoing phase transformations. In the framework of generalized standard solids due to Halphen and Nguyen (see [HaN75]), the unknowns are the displacement field $u$ and an internal variable $z$ and the problem is described by the momentum equilibrium equation combined with a flow rule for the evolution of the internal variable. A very powerful tool to study such problems is the so-called energetic formulation, introduced in [MT04] [Mic05] and later on developed and intensively applied in [FrM06] [MR06] [MR07] [Mic07] [MP07] [MR08]. Note that coupling rate-independent processes with rate-dependent processes makes, in general, the problems much more difficult; see for instance [ELM06] [MPM08] [BaR08] [Rou09a] [Rou09b].

In this paper, we are interested in coupling the rate-independent process with the thermal process, which is not rate-independent, and viscous damping. The model is based on the Helmholtz free energy $W(e(u), z, \nabla z, \theta)$, depending on the infinitesimal strain tensor $e(u) \overset{\text{def}}{=} \frac{1}{2}(\nabla u + \nabla u^T)$ for the displacement $u$, where $(\cdot)^T$ denotes the transpose of the tensor, the internal variable $z$ and the temperature $\theta$. For simplicity, we will omit any dependence on the material point $x \in \Omega$ and $t \in [0, T]$ with $T > 0$. We assume that $W$ can be decomposed as follows

$$W(e(u), z, \nabla z, \theta) \overset{\text{def}}{=} W_1(e(u), z, \nabla z) - W_0(\theta) + \theta W_2(e(u), z). \tag{1.1}$$

This partially linearized decomposition ensures that entropy separates the thermal and mechanical variables (see [23]). Let us emphasize that the last term $\theta W_2(e(u), z)$ allows for coupling effects between the temperature and the internal variable, which is motivated by the phenomenological models for shape memory alloys presented in Section 2. Since coupling terms will appear both in the momentum equilibrium equation and in the inclusion describing the evolution of the internal variable, this setting is more

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general than the one presented in \cite{Rou10}. We make the assumptions of small deformations. The problem is thus described by the following system

\begin{align}
&-\text{div}(\sigma_{el}+L\varepsilon(\dot{u})) = \ell, \quad \sigma_{el} \overset{\text{def}}{=} D_{e(u)}W(e(u), z, \nabla z, \theta), \\
&\partial \Psi(\dot{z}) + M\dot{z} + \sigma_{in} \geq 0, \quad \sigma_{in} \overset{\text{def}}{=} D_{z}W(e(u), z, \nabla z, \theta) - \text{div} D_{\nabla z}W(e(u), z, \nabla z, \theta), \\
&c(\theta)\dot{\theta} - \text{div}(\kappa e(u), z, \theta)\nabla \theta = L\varepsilon(\dot{u}) : e(\dot{u}) + \theta D_{\nabla z}W_{2}(e(u), z) + \Psi(\dot{z}) + M\dot{z}:\dot{z}.
\end{align}

Here $\Psi$ denotes the dissipation potential. As it is common in modeling hysteresis effect in mechanics, we assume that $\Psi$ is positively homogeneous of degree 1, i.e., $\Psi(\lambda v) = \lambda \Psi(v)$ for all $\lambda \geq 0$. The viscosity tensors are denoted $L$ and $M$, $c(\theta)$ is the heat capacity and $\kappa(e(u), z, \theta)$ is the conductivity. As usual, $(\cdot)^{i}, D_{i}^{(\cdot)}$ and $\partial_{i}$ denote the time derivative $\frac{\partial}{\partial t}$, the $i$-th derivative with respect to $z$ and the subdifferential in the sense of convex analysis (for more details see \cite{Bre73}), respectively. Observe that (1.2a), (1.2b) and (1.2c) are usually called the momentum equilibrium equation, the flow rule and the heat-transfer equation, respectively.

The paper is organized as follows. In Section 2, the thermodynamic consistency is justified and some illustrative examples are presented. Then the mathematical formulation of the problem in terms of displacement, internal variables and temperature is presented in Section 3. Our problem is reformulated in terms of enthalpy, which is a crucial ingredient to prove the existence result. Sections 4, 5 and 6 are devoted to the proof of a local existence result by a fixed-point argument. More precisely, in Section 4 we consider first the system composed by the momentum equilibrium equation and flow rule for a given temperature $\theta$ and we prove existence and regularity results. Next in Section 5 we recall existence and regularity results for the enthalpy equation for any given right hand side. Then a local existence result follows in Section 6 by using a fixed-point argument. Finally a global energy estimate is established in Section 7 leading to a global existence result for the system (1.2).

## 2 Mechanical model

We justify here the thermodynamic consistency of the model (1.2). Starting from the Helmholtz free energy $W$, we introduce the specific entropy $s$ via the Gibb’s relation

\begin{equation}
s \overset{\text{def}}{=} -D_{\theta}W(e(u), z, \nabla z, \theta),
\end{equation}

and the internal energy

\begin{equation}
W_{in}(e(u), z, \nabla z, \theta) \overset{\text{def}}{=} W(e(u), z, \nabla z, \theta) + \theta s.
\end{equation}

Then the entropy equation is given by

\begin{equation}
\theta \dot{s} + \text{div}(j) = \xi,
\end{equation}

where $j$ is the heat flux and $\xi$ is the dissipation rate. We get

\begin{equation}
\xi = L\varepsilon(\dot{u}) : e(\dot{u}) + M\dot{z}:\dot{z} + \Psi(\dot{z}) \geq 0,
\end{equation}

and, assuming Fourier’s law for the temperature, we have

\begin{equation}
 j = -\kappa(e(u), z, \theta)\nabla \theta.
\end{equation}

We can check now that the second law of thermodynamics is satisfied if $\theta > 0$. Indeed, assuming that the system is thermally isolated, we may divide (2.3) by $\theta$ and Green’s formula yields

\begin{align}
\int_{\Omega} \dot{s} \, dx &= \int_{\Omega} \frac{\text{div}(\kappa(e(u), z, \theta)\nabla \theta)}{\theta} \, dx + \int_{\Omega} \frac{L\varepsilon(\dot{u}) : e(\dot{u}) + M\dot{z}:\dot{z} + \Psi(\dot{z})}{\theta} \, dx \\
&= \int_{\Omega} \frac{\kappa(e(u), z, \theta)\nabla \theta}{\theta^{2}} \, dx + \int_{\Omega} \frac{L\varepsilon(\dot{u}) : e(\dot{u}) + M\dot{z}:\dot{z} + \Psi(\dot{z})}{\theta} \, dx \geq 0.
\end{align}

We differentiate now $W_{in}(e(u), z, \nabla z, \theta)$ with respect to time, we obtain by using the chain rule and (2.2) that

\begin{align}
\dot{W}_{in}(e(u), z, \nabla z, \theta) &= D_{e(u)}W(e(u), z, \nabla z, \theta) : e(\dot{u}) + D_{z}W(e(u), z, \nabla z, \theta) : \dot{z} + D_{\nabla z}W(e(u), z, \nabla z, \theta) : \nabla \dot{z} + \theta \dot{s}.
\end{align}
We integrate (2.4) over \( \Omega \), thus we use the Green’s formula and (2.3), we find

\[
\int_{\Omega} \dot{W}_{\text{in}}(e(u), z, \nabla z, \theta) \, dx = \int_{\Omega} D_{\text{e}}(u) W(e(u), z, \nabla z, \theta) : c(\dot{u}) \, dx + \int_{\Omega} D_{\text{z}} W(e(u), z, \nabla z, \theta) : \dot{z} \, dx \\
+ \int_{\Omega} D_{\text{vz}} W(e(u), z, \nabla z, \theta) \cdot \nabla \dot{z} \, dx + \int_{\Omega} (\nabla(\kappa(e(u), z, \theta) \nabla \theta) + \text{Le}(\dot{u}) : c(\dot{u}) + M \dot{z} : \dot{z} + \Psi(\dot{z})) \, dx.
\]  

(2.5)

On the one hand, we multiply (1.2b) by \( \dot{u} \), and we integrate this expression over \( \Omega \) to get

\[
\int_{\Omega} D_{\text{e}}(u) W(e(u), z, \nabla z, \theta) : c(\dot{u}) \, dx + \int_{\Omega} \text{Le}(\dot{u}) : c(\dot{u}) \, dx = \int_{\Omega} \ell \cdot \dot{u} \, dx.
\]  

(2.6)

On the other hand, the definition of the subdifferential \( \partial \Psi(\dot{z}) \) leads to the variational equality associated to (1.2b)

\[
\int_{\Omega} D_{\text{z}} W(e(u), z, \nabla z, \theta) : \dot{z} \, dx + \int_{\Omega} D_{\text{vz}} W(e(u), z, \nabla z, \theta) \cdot \nabla \dot{z} \, dx + \int_{\Omega} M \dot{z} : \dot{z} \, dx + \int_{\Omega} \Psi(\dot{z}) \, dx = 0.
\]  

(2.7)

We use (2.6) and (2.7) into (2.5), we obtain

\[
\int_{\Omega} \dot{W}_{\text{in}}(e(u), z, \nabla z, \theta) \, dx = \int_{\Omega} \ell \cdot \dot{u} \, dx + \int_{\partial \Omega} \kappa(e(u), z, \theta) \nabla \theta : \eta \, dx.
\]

This means that the total energy balance can be expressed in terms of the internal energy, which is the sum of power of external load and heat.

Note that from (2.1), we get

\[
s \overset{\text{def}}{=} D_\theta W_0(\theta) - W_2(e(u), z),
\]

(2.8)

and

\[
W_{\text{in}}(e(u), z, \nabla z, \theta) \overset{\text{def}}{=} W_1(e(u), z, \nabla z, \theta) + \theta D_\theta W_0(\theta) - W_0(\theta).
\]

(2.9)

We use (2.9) into (2.4), we may deduce the heat-transfer equation (1.2b) with the heat capacity given by \( c(\theta) = \theta D_\theta^2 W_0(\theta) \).

Motivated by the study of shape memory alloys, we will focus in the rest of the paper on the special case where \( W_1 \) is given by

\[
W_1(e(u), z, \nabla z) \overset{\text{def}}{=} \frac{1}{2} \mathbb{E}(e(u) - z) : (e(u) - z) + \frac{\nu}{2} |\nabla z|^2 + H_1(z),
\]

where the internal variable \( z \) is a deviatoric \( d \times d \) tensor, \( H_1 \) is a hardening functional and the term \( \frac{\nu}{2} |\nabla z|^2 \), \( \nu > 0 \), takes into account some nonlocal interaction effect for the internal variable. Moreover we will assume

\[
W_2(e(u), z) \overset{\text{def}}{=} \alpha \text{tr}(e(u)) + H_2(z).
\]

Here \( \alpha I \) with \( \alpha \geq 0 \) and I the identity matrix, is the isothermal expansion tensor. Furthermore, the sum \( H(z, \theta) \overset{\text{def}}{=} H_1(z) + \theta H_2(z) \) may be interpreted as a temperature-dependent hardening functional. Note that the system (1.2) is then rewritten as

\[
- \text{div}(\mathbb{E}(e(u) - z) + \alpha \theta I + \text{Le}(\dot{u})) = \ell,
\]

\[
\partial \Psi(\dot{z}) + M \dot{z} - \mathbb{E}(e(u) - z) + D_{\dot{z}} H_1(z) + \theta D_{\dot{z}} H_2(z) - \nu \Delta z \ni 0,
\]

\[
c(\theta) \dot{\theta} - \text{div}(\kappa(e(u), z, \theta) \nabla \theta) = \text{Le}(\dot{u}) : c(\dot{u}) + \theta (\alpha \text{tr}(e(u)) + D_{\dot{z}} H_2(z) : \dot{z}) + \Psi(\dot{z}) + M \dot{z} : \dot{z}.
\]

We may illustrate the presented model by a nontrivial example, namely, a three-dimensional macroscopic phenomenological model for shape-memory polycrystalline materials undergoing phase transformations. This model was introduced by Souza et al (SMZ98) and by Auricchio et al (see [AuP02]). The internal variable describes the inelastic part of the deformation due to the martensitic phase transformation and the hardening functional \( H \) takes the form

\[
H(z, \theta) \overset{\text{def}}{=} c_1(\theta) |z| + c_2(\theta) |z|^2 + \chi(z).
\]
Here $\chi : \mathbb{R}^{d_{\text{dev}}} \rightarrow [0, +\infty]$ denotes the indicator function of the ball $\{z \in \mathbb{R}^{d_{\text{dev}}} : |z| \leq c_3(\theta)\}$ and the coefficients $c_i(\theta)$ are positive real numbers. Let us observe that $c_1(\theta) > 0$ is an activation threshold for initiation of martensitic phase transformations, $c_2(\theta)$ measures the occurrence of hardening with respect to the internal variable $z$ and $c_3(\theta)$ represents the maximum modulus of transformation strain that can be obtained by alignment of martensitic variants.

The dependence of the coefficients $c_i$, $i = 1, 2, 3$, with respect to $\theta$ is due to a strong thermo-mechanical constitutive coupling coming from the latent heat absorption or release, which is one of the main features of the behavior of shape-memory alloys (see [AuP04] for more details).

For mathematical purposes, we should regularize the hardening functional as in [MiP07]. Namely, we replace $H(z, \theta)$ by

$$H^\delta(z, \theta) \overset{\text{def}}{=} c_1(\theta) \sqrt{\delta^2 + |z|^2} + c_2(\theta)|z|^2 + \frac{(|z| - e_3(\theta))^4}{\delta(1 + |z|^2)},$$

where $\delta > 0$ is a small parameter. Then, assuming that the mappings $c_i$, $i = 1, 2, 3$, are of class $C^1$ we can consider an affine approximation of $H^\delta(z, \theta)$ as $H^\delta_1(z) + \theta H^\delta_2(z)$ with

$$H^\delta_1(z) = c_1 \sqrt{\delta^2 + |z|^2} + c_2 |z|^2 + \frac{(|z| - e_3)^4}{\delta(1 + |z|^2)},$$

where $c_i > 0$, $i = 1, 2, 3$.

Let us emphasize that the existence proof presented in the next sections can be easily extended to more general models of shape memory alloys for which $W_1$ is given by

$$W_1(e(u), z, \nabla z) \overset{\text{def}}{=} \frac{1}{2} \mathbb{E}(e(u) - E(z)) : (e(u) - E(z)) + \frac{\nu}{2} |\nabla z|^2 + H_1(z),$$

where the internal variable $z$ is a vector of $\mathbb{R}^{N-1}$, with $N \geq 2$, and $E$ is an affine mapping from $\mathbb{R}^{N-1}$ to the set of $d \times d$ deviatoric tensors. In such models, $N$ is the total number of phases, i.e., the austenite and all the variants of martensite, and the components $z_1, \ldots, z_{N-1}$ of $z$ and $z_N \overset{\text{def}}{=} 1 - \sum_{k=1}^{N-1} z_k$ are interpreted as phase fractions. Then $E(z)$ is the effective transformation strain of the mixture, given by

$$E(z) \overset{\text{def}}{=} \sum_{k=1}^{N-1} z_k E_k + \left(1 - \sum_{k=1}^{N-1} z_k\right) E_N,$$

where $E_k$ is the transformation strain of the phase $k$ and the temperature dependent hardening functional $H(z, \theta)$ is the sum of a smooth part $w(z, \theta)$ and the indicator function of the set $[0, 1]^{N-1}$ (see [MiT99, Mic00, HacG02, GMH02, MTL02, GH07]). The hardening functional can be regularized in a similar way as in the previous example by replacing $H(z, \theta)$ by

$$H^\delta(z, \theta) = w(z, \theta) + \sum_{k=1}^{N-1} \frac{((z_k)^4 + ((z_k-1)^4)}{\delta(1 + |z_k|^2)},$$

and we may consider an affine approximation of $w(z, \theta)$ as $w_1(z) + \theta w_2(z)$. For more details on this example the reader is referred to [PaP11].

3 Mathematical formulation

We consider a reference configuration $\Omega \subset \mathbb{R}^3$. We assume that $\Omega$ is an bounded domain such that $\partial \Omega$ is of class $C^{2+\rho}$. We will denote by $\mathbb{R}^{3 \times 3}_{\text{sym}}$ the space of symmetric $3 \times 3$ tensors endowed with the natural scalar product $v:w \overset{\text{def}}{=} \text{tr}(v^T w)$ and the corresponding norm $|v|^2 \overset{\text{def}}{=} v:v$ for all $v, w \in \mathbb{R}^{3 \times 3}_{\text{sym}}$. In particular, we assume that

$$W_1(e(u), z, \nabla z) \overset{\text{def}}{=} \frac{1}{2} \mathbb{E}(e(u) - z) : (e(u) - z) + \frac{\nu}{2} |\nabla z|^2 + H_1(z) \quad \text{and} \quad W_2(e(u), z) \overset{\text{def}}{=} \alpha \text{tr}(e(u)) + H_2(z),$$

where $\nu > 0$, $\alpha \geq 0$, is the isotropic thermal expansion coefficient, $\mathbb{E}$ denotes the elastic tensor and $H_i$, $i = 1, 2$, two hardening functionals. Given a function $\ell : \Omega \times (0, T) \rightarrow \mathbb{R}^3$, we look for a displacement $u$: 
Global existence result for phase transformations with heat transfer in shape memory alloys

\[ \Omega \times (0, T) \rightarrow \mathbb{R}^3 \]

a matrix of internal variables \( z : \Omega \times (0, T) \rightarrow \mathbb{R}^{3 \times 3} \) and a temperature \( \theta : \Omega \times (0, T) \rightarrow \mathbb{R} \)
satisfying the following system:

\[
\begin{align*}
- \text{div} (E(e(u) - z) + \alpha \theta I + L e(\dot{u})) &= \ell, \\
\partial \Psi (\dot{z}) + M \dot{z} - E(e(u) - z) + D_z H_i(z) + \theta D_z H_2(z) - \nu \Delta z &\geq 0, \\
c(\theta) \dot{\theta} - \text{div} (\kappa c(e(u), z, \theta) \nabla \theta) &= L e(\dot{u}) : e(\dot{u}) + \theta (\alpha \text{tr}(e(u))) + D_z H_2(z) : \dot{z} + \Psi (\dot{z}) + M \dot{z} : \dot{z}.
\end{align*}
\]

We have naturally to prescribe initial conditions for the displacement, the internal variables, and the
temperature, namely

\[
u(\cdot, 0) = u^0, \quad z(\cdot, 0) = z^0, \quad \theta(\cdot, 0) = \theta^0.
\]

The problem is to be completed with boundary conditions. More precisely, we suppose here that

\[
\begin{align*}
u_{1, \alpha \alpha} &= 0, \quad \nabla z \cdot \eta_{1, \alpha} = 0, \quad \kappa \nabla \eta \cdot \eta_{1, \alpha} = 0,
\end{align*}
\]

where \( \eta \) denotes the outward normal to the boundary \( \partial \Omega \) of \( \Omega \). The original problem \( \text{(3.1)} \) can be rewritten in terms of enthalpy instead of temperature by employing the so-called enthalpy transformation

\[
g(\theta) = \dot{\theta} \overset{\text{def}}{=} \int_0^\theta c(s) \, ds.
\]

Clearly, \( g \) is the unique primitive of the function \( c \), which is supposed to be continuous, such that \( g(0) = 0 \). Furthermore, we will assume that for all \( s \geq 0 \), \( c(s) \geq c^c > 0 \) where \( c^c \) is a constant. Hence we deduce that \( g \) is a bijection from \([0, \infty)\) into \([0, \infty)\). We define

\[
\zeta(\dot{\theta}) \overset{\text{def}}{=} \begin{cases} g^{-1}(\dot{\theta}) & \text{if } \dot{\theta} \geq 0, \\ 0 & \text{otherwise}, \end{cases}
\]

\[
\kappa^c (e(u), z, \theta) \overset{\text{def}}{=} \frac{\kappa c(e(u), z, \zeta(\dot{\theta}))}{c(\zeta(\dot{\theta}))},
\]

where \( g^{-1} \) is the inverse of \( g \). For more details on the enthalpy transformation, the reader is referred to

\[ \text{([Rou09])} \]

and the references therein. Therefore the system \( \text{(3.1)} \) is transformed into the following form

\[
\begin{align*}
- \text{div} (E(e(u) - z) + \alpha \zeta(\dot{\theta}) I + L e(\dot{u})) &= \ell, \\
\partial \Psi (\dot{z}) + M \dot{z} - E(e(u) - z) + D_z H_i(z) + \zeta(\dot{\theta}) D_z H_2(z) - \nu \Delta z &\geq 0, \\
\dot{\theta} - \text{div} (\kappa^c (e(u), z, \dot{\theta}) \nabla \theta) &= L e(\dot{u}) : e(\dot{u}) + \zeta(\dot{\theta}) (\alpha \text{tr}(e(u))) + D_z H_2(z) : \dot{z} + \Psi (\dot{z}) + M \dot{z} : \dot{z},
\end{align*}
\]

with boundary conditions

\[
\begin{align*}
u_{1, \alpha \alpha} &= 0, \quad \nabla z \cdot \eta_{1, \alpha} = 0, \quad \kappa^c \nabla \eta \cdot \eta_{1, \alpha} = 0,
\end{align*}
\]

and initial conditions

\[
u(\cdot, 0) = u^0, \quad z(\cdot, 0) = z^0, \quad \theta(\cdot, 0) = \theta^0 = g(\theta^0).
\]

The identity \( \text{(3.6c)} \) is called the enthalpy equation. As usual Korn’s inequality will play a role in the
mathematical analysis developed in the next sections. We have assumed that \( \partial \Omega \) is of class \( C^{2+\rho} \), so that
the Korn’s inequality holds, i.e.

\[
\exists C_{\text{Korn}} > 0 \ \forall u \in H^1_0(\Omega) : \| e(u) \|^2_{L^2(\Omega)} \geq C_{\text{Korn}} \| u \|^2_{H^1(\Omega)},
\]

(see \[ \text{[KoOSS, DIL76]} \]).

Let us introduce now the assumptions on the dissipation potential \( \Psi \), on the hardening functions \( H_i \),
\( i = 1, 2 \), and on the data \( E, L, M, \ell, c \overset{\text{def}}{=} c(\theta) \) and \( \kappa \overset{\text{def}}{=} \kappa (e(u), z, \theta) \), which will allow us to obtain some
regularity properties that are needed to prove the existence result.

We assume that the dissipation potential \( \Psi \) is positively homogeneous of degree 1 and satisfies the
triangle inequality, namely, we have

\[
\begin{align*}
\forall \zeta \geq 0 \ \forall z \in \mathbb{R}^{3 \times 3}_{\text{sym}} : \Psi (\gamma z) &= \gamma \Psi (z), \\
\exists C^\Psi > 0 \ \forall z \in \mathbb{R}^{3 \times 3}_{\text{sym}} : 0 \leq \Psi (z) \leq C^\Psi |z|, \\
\forall z_1, z_2 \in \mathbb{R}^{3 \times 3}_{\text{sym}} : \Psi (z_1 + z_2) \leq \Psi (z_1) + \Psi (z_2).
\end{align*}
\]
It is clear that (3.10a), (3.10b) and (3.10c) imply that $\Psi$ is convex and continuous. We impose that the hardening functionals $H_i$, $i = 1, 2$, belong to $C^2(\mathbb{R}^{3\times 3}; \mathbb{R})$ and satisfy the following inequalities
\begin{align}
\exists c_{H_i} : 0 < c_{H_i} & \forall z \in \mathbb{R}^{3\times 3} : 1_H(z) \geq c_{H_i} |z|^2 - c_{H_i}, \quad (3.11a) \\
\exists c_{zz_i} & > 0 \forall z \in \mathbb{R}^{3\times 3} : |D^2_H(z)| \leq c_{zz_i}. \quad (3.11b)
\end{align}

Note that (3.11b) leads to
\begin{align}
\exists C_{H_i} > 0 \forall z \in \mathbb{R}^{3\times 3}_{\text{dev}} : |D_z H_i(z)| \leq C_{H_i}(1 + |z|), \quad |H_i(z)| \leq C_{H_i}(1 + |z|^2). \quad (3.12)
\end{align}

The elastic tensor $E : \Omega \to \mathbb{L}(\mathbb{R}^{3\times 3}, \mathbb{R}^{3\times 3})$ is a symmetric positive definite operator such that
\begin{align}
\exists \kappa > 0 & \forall z \in L^2(\Omega; \mathbb{R}^{3\times 3}) : \kappa \leq \int \Omega |D_z z|^2 \leq \int \Omega |z|^2 \leq 2, \quad (3.13a) \\
\exists \lambda & > 0 \forall i, j, k = 1, 2, 3 : \lambda \in (0, \infty). \quad (3.13b)
\end{align}

We assume that $L$ and $M$ are symmetric positive definite tensors. This implies that
\begin{align}
\exists \lambda^L, \lambda^M & > 0 \forall z \in \mathbb{R}^{3\times 3} : \lambda^L |z|^2 \leq L_z z \leq \lambda^M |z|^2, \quad (3.14a) \\
\exists \mu^L, \mu^M & > 0 \forall z \in \mathbb{R}^{3\times 3} : \mu^L |z|^2 \leq M_z z \leq \mu^M |z|^2. \quad (3.14b)
\end{align}

We consider that $\ell$ is an external loading satisfying
\begin{align}
\ell & \in H^1(0, T; L^2(\Omega)). \quad (3.15)
\end{align}

Finally, for the heat capacity $c$ and the conductivity $\kappa^c$ we assume that
\begin{align}
c : [0, \infty) & \to [0, \infty) is continuous, \quad (3.16a) \\
\exists \beta_1 & \geq 2 \exists \beta > 0 : c^c(1 + \theta)^{\beta_1 - 1} \leq c(\theta), \quad (3.16b) \\
\kappa^c : \mathbb{R}^{3\times 3} & \times \mathbb{R}^{3\times 3} \times \mathbb{R} \to \mathbb{R}^{3\times 3} \text{ is continuous, bounded and uniformly positive definite, i.e} \quad (3.16c) \\
\exists \kappa^c & > 0 \forall (e, z, \vartheta) \in \mathbb{R}^{3\times 3}_{\text{sym}} \times \mathbb{R}^{3\times 3} \times \mathbb{R} : \kappa^c(e, z, \vartheta) |v|^2 \leq \kappa^c(e, z, \vartheta)|v|^2, \quad (3.16d) \\
\exists C^c & > 0 \forall (e, z, \vartheta) \in \mathbb{R}^{3\times 3}_{\text{sym}} \times \mathbb{R}^{3\times 3} \times \mathbb{R} : |\kappa^c(e, z, \vartheta)| \leq C^c. \quad (3.16e)
\end{align}

Let us emphasize that the three dimensional model for shape memory alloys presented in Section 2 fulfills the previous assumptions so that we may apply the abstract result obtained in the next sections.

Let us end this section by some comments about the proof strategy. In order to obtain a local existence result for the coupled problem (3.6)–(3.8), a fixed point argument will be used. More precisely, for any given $\tilde{\vartheta}$, we define $\theta = \zeta(\tilde{\vartheta})$ and we solve first the system composed by the momentum equilibrium equation and the flow rule (3.11)–(3.11), then we solve the enthalpy equation (3.6c) with $\kappa^c = \kappa^c(e(u), z, \zeta(\tilde{\vartheta}))$. This allows us to define a mapping
\begin{equation}
\phi : \tilde{\vartheta} \mapsto \vartheta,
\end{equation}
and our aim is to prove that this mapping satisfies the assumptions of Schauder’s fixed point theorem. Therefore, let us consider a given $\tilde{\vartheta} \in L^p(0, T; L^p(\Omega))$ with $\tilde{p} \geq 1$ and $\bar{p} \geq 1$. We define $\theta = \zeta(\tilde{\vartheta})$. Since $\zeta$ is a Lipschitz continuous mapping from $\mathbb{R}$ to $\mathbb{R}$, we infer that the mapping
\begin{equation}
\phi_1 : L^\tilde{p}(0, T; L^\tilde{p}(\Omega)) \to L^\tilde{p}(0, T; L^\tilde{p}(\Omega)) \quad (3.17)
\end{equation}
is also Lipschitz continuous. Furthermore (3.17b) implies that
\begin{equation}
\forall \theta \in [0, \infty) : \frac{\kappa^c(1 + \theta)^{\beta_1 - 1}}{\lambda^c} g_1(\theta) \leq g(\theta). \quad (3.18)
\end{equation}
Thus we have
\[ \forall \theta \in [0, \infty) : 0 \leq \zeta(\theta) \leq \zeta_1(\theta) \overset{\text{def}}{=} g_1^{-1}(\theta), \]
and
\[ \forall \theta \in \mathbb{R} : |\zeta(\theta)| \leq \left( \frac{\theta^+}{\max(\theta, 0)} + 1 \right)^{\frac{1}{q}} - 1, \]
with \( \theta^+ \overset{\text{def}}{=} \max(\theta, 0) \) for all \( \theta \in \mathbb{R} \). It follows that
\[ \forall \beta \in [1, \beta_1] \forall \theta \in \mathbb{R} : |\zeta(\theta)| \leq \left( \frac{\beta^+}{\max(\theta, 0)} + 1 \right)^{\frac{1}{q}} - 1 \leq \left( \frac{\theta^+}{\max(\theta, 0)} + 1 \right)^{\frac{1}{q}}. \] (3.18)
Hence, for all \( \beta \in [1, \beta_1] \) and for all \( \theta \in L^{q}(0,T;L^{p}(\Omega)) \), we have \( \theta = \zeta(\theta) \in L^{p}((0,T;L^{p}(\Omega)) \) with
\[ \|\theta\|_{L^{p}(0,T;L^{p}(\Omega))} \leq \left( \frac{\theta^+}{\max(\theta, 0)} + 1 \right)^{\frac{1}{q}} \|\theta\|_{L^{p}(0,T;L^{p}(\Omega))}. \]

In the rest of the paper, we will assume that \( q > 4 \) and \( p = 2 \).

When there is not any confusion, we will use simply the notation \( X(\Omega) \) instead of \( X(\Omega; Y) \) where \( X \) is a functional space and \( Y \) is a vectorial space.

4 Existence and regularity results for the system composed by the momentum equilibrium equation and the flow rule

This section is devoted to the proof of existence and uniqueness results for the system composed by the momentum equilibrium equation and the flow rule (3.13–3.14) under the consideration that \( \theta = \zeta(\theta) \) is given in a bounded subset of \( L^{q}(0,T;L^{p}(\Omega)) \) with \( q = \beta_1 q \) and \( p \in [4, \min(\beta_1 p, 6)] \). More precisely, we look for a solution of the problem (P_u):
\[- \text{div}(\mathbb{E}(e(u) - z) + \alpha \mathbb{I} + \mathbb{E}(e(u))) = \ell, \]
\[ \partial \Psi(z) + M \dot{z} - \mathbb{E}(e(u) - z) + D_2 H_1(z) + \theta D_2 H_2(z) - \nu \Delta z \ni 0, \] (4.1a)
with initial conditions
\[ u(\cdot, 0) = u^0, \quad z(\cdot, 0) = z^0, \] (4.2)
and boundary conditions
\[ u|_{\partial \Omega} = 0, \quad \nabla z \cdot n|_{\partial \Omega} = 0. \] (4.3)
Furthermore we will establish some a priori estimates and some regularity results for the solution of (P_u).

As a first step, we use classical results for Partial Differential Equations (PDE) and Ordinary Differential Equations (ODE) to obtain an existence result. Let \( A : H^1(\Omega) \to (H^1(\Omega))' \) be the linear continuous mapping defined by
\[ \forall (u, v) \in (H^1(\Omega))^2 : \langle Au, v \rangle_{(H^1(\Omega))', H^1(\Omega)} \overset{\text{def}}{=} \int_{\Omega} \nu M^{-1} \nabla u : \nabla v \, dx. \]
Classical results about elliptic operators implies that \( A \) generates an analytic semigroup on \( L^2(\Omega) \), which extends to a \( C^0 \)-semigroup of contractions on \( L^p(\Omega) \). We denote by \( A_p \) (resp. \( A_{\xi} \)) the realization of its generator in \( L^p(\Omega) \) (resp. \( L^{p/2}(\Omega) \)) and by \( X_{q,p}(\Omega) \) the intersection of interpolation spaces
\[ X_{q,p}(\Omega) \overset{\text{def}}{=} (L^p(\Omega), D(A_p))_{1-\frac{2}{q}, \frac{2}{q}} \cap (L^{p/2}(\Omega), D(A_{\xi}))_{1-\frac{4}{q}, \frac{4}{q}}, \]
(see for instance [HiR08, PrS01] and the references therein). Here \( D(A_p) \) (resp. \( D(A_{\xi}) \)) denotes the domain of \( A_p \) (resp. \( A_{\xi} \)).

In the sequel, the notations for the constants introduced in the proofs are valid only in the proof and we also use the set \( Q_{\tau} \overset{\text{def}}{=} \Omega \times (0, \tau) \) with \( \tau \in [0, T] \).

**Theorem 4.1 (Existence for (P_u))** Let \( \theta \) be given in \( L^{q}(0,T;L^{p}(\Omega)) \). Assume that \( u^0 \in H^1(\Omega) \), \( z^0 \in \text{H}^1(\Omega) \) and that (3.10), (3.11), (3.14) and (3.16) hold. Then the problem (4.1)–(4.3) admits a solution \( (u, z) \in H^1(\Omega) \times L^2(\Omega) \cap L^\infty(0,T;H^1_0(\Omega) \times H^1(\Omega)) \). Furthermore if \( z^0 \in X_{q,p}(\Omega) \) then \( z \in L^2(0,T;H^2(\Omega)) \cap C^0([0,T];H^1(\Omega)) \) and \( \dot{z} \in L^2(0,T;L^2(\Omega)) \).
Proof. Observe first that for all \( f \in L^2(0,T; (H^1_0(\Omega); \mathbb{R}^3)') \) and for all \( u^* \in H^1_0(\Omega; \mathbb{R}^3) \) the following problem
\[
- \text{div}(\mathcal{E}u(\cdot)+\mathcal{L}(\dot{u}(\cdot))) = f,
\]
with initial conditions
\[
u \cdot 0 = u^* \in H^1_0(\Omega; \mathbb{R}^3),
\]
and boundary conditions
\[
u |_{\partial \Omega} = 0,
\]
possesses a unique solution \( u^* \defeq \mathcal{L}(u^*, f) \in H^1(0,T; (H^1_0(\Omega)) \cap C^0([0,T]; H^1_0(\Omega))). \) Note that
\[
\forall f_1, f_2 \in L^2(0,T; (H^1_0(\Omega); \mathbb{R}^3)') \ \forall u^* \in H^1_0(\Omega; \mathbb{R}^3) : \ \mathcal{L}(u^*, f_1 + f_2) = \mathcal{L}(u^*, f_1) + \mathcal{L}(0, f_2),
\]
and the mapping
\[
\mathcal{L}_0(\cdot) \defeq \mathcal{L}(0, \cdot) : L^2(0,T; (H^1_0(\Omega)')) \rightarrow H^1(0,T; (H^1_0(\Omega)) \cap C^0([0,T]; H^1_0(\Omega)),
\]
\[
f \mapsto u,
\]
is linear and continuous. Moreover, the classical energy estimate gives
\[
\forall t \in [0,T] : c^2 \|e(u(\cdot,t))\|_{L^2(\Omega)}^2 + c^3 \|\dot{u}(\cdot, s)\|_{L^1(\Omega)}^2 ds \leq \frac{1}{c^3 Korn} \int_0^t \|f(\cdot, s)\|_{[H^1_0(\Omega)]'}^2 ds,
\]
for all \( f \in L^2(0,T; (H^1_0(\Omega)')) \) and \( u = \mathcal{L}_0(f). \) It follows that \((4.1a)\) and \((4.1b)\) can be rewritten as
\[
\partial \Psi(\dot{z}) + \mathbb{M} \dot{z} + \mathbb{E} z + D_z H_1(z) + \theta D_z H_2(z) - \nu \Delta z + g_1(\theta) + g_2(z) \geq 0,
\]
with initial conditions
\[
z(\cdot, 0) = z^0 \in H^1(\Omega; \mathbb{R}^{3 \times 3}),
\]
and boundary conditions
\[
\nabla z \cdot \eta |_{\partial \Omega} = 0,
\]
where
\[
g_1(\theta) \defeq -\mathcal{E}(\mathcal{L}(u^0, \ell)) - \mathcal{E}(\mathcal{L}_0(\text{div}(\alpha \theta I))) \in H^1(0,T; L^2(\Omega; \mathbb{R}^{3 \times 3}) \text{ and})
\]
\[
g_2 : L^2(0,T; L^2(\Omega; \mathbb{R}^{3 \times 3})) \rightarrow H^1(0,T; L^2(\Omega; \mathbb{R}^{3 \times 3})),
\]
\[
z \mapsto \mathcal{E}(\mathcal{L}_0(\text{div}(\mathbb{E} z))).
\]
Let \( \varphi : L^2(\Omega; \mathbb{R}^{3 \times 3}) \rightarrow (-\infty, +\infty] \) defined by
\[
z \mapsto \begin{cases} \frac{\varphi}{2} \|\nabla z\|^2_{L^2(\Omega)} & \text{if } z \in H^1(\Omega; \mathbb{R}^{3 \times 3}), \\ +\infty & \text{otherwise,} \end{cases}
\]
is a proper, convex lower semicontinuous function. Clearly, \( \partial \varphi \) is a maximal monotone operator, for more details on the maximal monotone operators and their properties, the reader is referred to [Bre73]. The resolvent of the subdifferential \( \partial \varphi \) is defined by
\[
\forall \epsilon > 0 : \mathcal{J}_\epsilon \defeq (1+\epsilon \partial \varphi)^{-1}.
\]
Note that the resolvent \( \mathcal{J}_\epsilon \) is a contraction defined on all \( L^2(\Omega; \mathbb{R}^{3 \times 3}) \). Let us introduce also
\[
\forall \epsilon > 0 \ \forall z \in L^2(\Omega; \mathbb{R}^{3 \times 3}) : \varphi_\epsilon(z) \defeq \min_{z \in L^2(\Omega; \mathbb{R}^{3 \times 3})} \left\{ \frac{1}{2\epsilon} \|z-\bar{z}\|^2_{L^2(\Omega)} + \varphi(\bar{z}) \right\},
\]
which is a convex and Fréchet differentiable mapping from \( L^2(\Omega; \mathbb{R}^{3 \times 3}) \) to \( \mathbb{R} \). Using [Bre73] Prop. 2.11, we know that the Yosida approximation of \( \partial \varphi \) coincide with the Fréchet differential of \( \varphi_\epsilon \), i.e., we have
\[
\forall \epsilon > 0 : \partial \varphi_\epsilon \overset{\text{def}}{=} \frac{1}{2}(1-\mathcal{J}_\epsilon),
\]
and \( \partial \varphi \) is a monotone and \( \frac{1}{2} \)-Lipschitz continuous mapping on \( L^2(\Omega; \mathbb{R}^{3 \times 3}) \).

We approximate now the problem \([3.3] - [4.7]\) by
\[
\partial \Psi(\tilde{z}_e) + M \tilde{z}_e + \partial \varphi_e(\tilde{z}_e) + E z_e + D_z H_1(z_e) \ni - (g_1(\theta_e) + g_2(J_e z_e) + \theta_e D_z H_2(J_e z_e)),
\]
with the initial condition
\[
z_e(0, \cdot) = z^0.
\]
Here \( \theta_e \in C^0_0(0, T) \otimes C^0_0(\Omega) \) and \( \partial \Psi(\tilde{z}_e) \) is taken in the sense of the \( L^2(\Omega) \)-extension of the subdifferential of the convex function \( \Psi \) (see examples 2.1.3 and 2.3.3 in \([Bre92]\)), i.e., \([4.9]\) is equivalent to
\[
\begin{align*}
&-M \tilde{z}_e(t, x) - \partial \varphi_e(z_e(t, x)) - E z_e(t, x) - D_z H_1(z_e(t, x)) - g_1(\theta_e(t, x)) \\
&- g_2(J_e z_e(t, x)) - \theta_e(t, x) D_z H_2(J_e z_e(t, x))) \in \partial \Psi(z_e(t, x)).
\end{align*}
\]
for all \( t \in [0, T] \) and almost every \( x \in \Omega \). Since we are looking for a solution with values in \( \mathbb{R}^{3 \times 3}_{\text{dev}} \), the test-functions should be taken in \( \mathbb{R}^{3 \times 3}_{\text{dev}} \). More precisely, we have
\[
\begin{align*}
\Psi(\tilde{z}) - \Psi(\tilde{z}(t, x)) + (M \tilde{z}_e(t, x) + \partial \varphi_e(z_e(t, x)) + E z_e(t, x) + D_z H_1(z_e(t, x)) \\
&+ g_1(\theta_e(t, x)) + g_2(J_e z_e(t, x)) + \theta_e(t, x) D_z H_2(J_e z_e(t, x)))(\tilde{z} - z_e(t, x)) \geq 0,
\end{align*}
\]
for all \( \tilde{z} \in \mathbb{R}^{3 \times 3}_{\text{dev}} \) and \( t \in [0, T] \) and almost every \( x \in \Omega \). Observe that \([4.11]\) is equivalent to
\[
\begin{align*}
\Psi_{\text{dev}}(\tilde{z}) - \Psi_{\text{dev}}(\tilde{z}(t, x)) + \text{Proj}_{\mathbb{R}^{3 \times 3}_{\text{dev}}}((M \tilde{z}_e(t, x) + \partial \varphi_e(z_e(t, x)) + E z_e(t, x) + D_z H_1(z_e(t, x)) \\
&+ g_1(\theta_e(t, x)) + g_2(J_e z_e(t, x)) + \theta_e(t, x) D_z H_2(J_e z_e(t, x)))(\tilde{z} - z_e(t, x)) \geq 0,
\end{align*}
\]
for all \( \tilde{z} \in \mathbb{R}^{3 \times 3}_{\text{dev}} \) and \( t \in [0, T] \) and almost every \( x \in \Omega \). Here \( \Psi_{\text{dev}} \) is the restriction of \( \Psi \) to \( \mathbb{R}^{3 \times 3}_{\text{dev}} \) and \( \text{Proj}_{\mathbb{R}^{3 \times 3}_{\text{dev}}} \) denotes the projection on \( \mathbb{R}^{3 \times 3}_{\text{dev}} \) relatively to the inner product of \( \mathbb{R}^{3 \times 3}_{\text{sym}} \). Using the definition of the subdifferential \( \partial \Psi_{\text{dev}}(z_e(t, x)) \) leads to the following differential inclusion
\[
\begin{align*}
&-\text{Proj}_{\mathbb{R}^{3 \times 3}_{\text{dev}}}((\partial \varphi_e(z_e(t, x)) + E z_e(t, x) + D_z H_1(z_e(t, x)) + g_1(\theta_e(t, x)) \\
&+ g_2(J_e z_e(t, x)) + \theta_e(t, x) D_z H_2(J_e z_e(t, x)))(\tilde{z} - z_e(t, x)) \ni \partial \Psi_{\text{dev}}(z_e(t, x)) + M \tilde{z}_e(t, x),
\end{align*}
\]
for all \( t \in [0, T] \) and almost every \( x \in \Omega \). Note that the operator \( \partial \Psi_{\text{dev}} + M \) is strongly monotone on \( \mathbb{R}^{3 \times 3}_{\text{dev}} \) as well as its \( L^2(\Omega) \)-extension, still denoted \( \partial \Psi_{\text{dev}} + M \), on \( L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}) \). This operator is invertible and its inverse is \( \frac{1}{2} \)-Lipschitz continuous. Thus \([4.9]\) can be rewritten as
\[
\begin{align*}
\dot{z}_e = \Phi(\theta_e, z_e),
\end{align*}
\]
with
\[
\Phi(\theta_e, z) \overset{\text{def}}{=} (\partial \Psi_{\text{dev}} + M)^{-1}(\text{Proj}_{\mathbb{R}^{3 \times 3}_{\text{dev}}}((\partial \varphi_e(z) - E z - D_z H_1(z) - g_1(\theta_e) - g_2(J_e z) - \theta_e D_z H_2(J_e z))),
\]
for all \( z \in L^2(0, T; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \) and we have to solve this differential equation for the unknown function \( z_e \). We may observe that we are not in the classical framework of ODE in \( L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}) \) (see \([Car90]\)), since \( g_2(J_e z) \) can not be defined for \( z \in L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}) \) but only for \( z \in L^2(0, T; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \). Nevertheless, we can solve \([4.12]\) with the Picard’s iterations technique. Indeed, \( \Phi(\theta_e, z) \in C^0([0, T]; L^2(\Omega)) \) for all \( z \in C^0([0, T]; L^2(\Omega)) \) and \( z_e \in C^1([0, T]; L^2(\Omega)) \) is a solution of \([4.9] - [4.10]\) if and only if \( z_e \) is a fixed point of the mapping
\[
\Lambda_e : C^0([0, T]; L^2(\Omega)) \rightarrow C^1([0, T]; L^2(\Omega))
\]
\[
z \mapsto \Lambda_e(z) : t \mapsto z^0 + \int_0^t \Phi(\theta_e(\cdot, s), z(\cdot, s)) \, ds.
\]
Let us assume that \( z_1, z_2 \in C^0([0, T]; L^2(\Omega)) \). By using \([3.11b]\) and \([3.13b]\), we find that
\[
\begin{align*}
&\|\Lambda_e(z_1(t, \cdot)) - \Lambda_e(z_2(t, \cdot))\|_{L^2(\Omega)} \leq \frac{1}{2} \int_0^t \left( \|\partial \varphi_e(z_1(\cdot, s)) - \partial \varphi_e(z_2(\cdot, s))\|_{L^2(\Omega)} + \|E\|_{L^\infty(\Omega)} + C^H_{z_2} + C^H_{z_2} \|\theta_e\|_{L^\infty(\Omega)}\right)\|z_1(\cdot, s) - z_2(\cdot, s)\|_{L^2(\Omega)} + \|g_2(J_e z_1(\cdot, s)) - g_2(J_e z_2(\cdot, s))\|_{L^2(\Omega)} \, ds,
\end{align*}
\]
for all $t \in [0, T]$. Since $\partial \varphi_x$ is $\frac{1}{\varepsilon}$-Lipschitz continuous on $L^2(\Omega)$ and $g_2$ is linear, it follows that

$$\forall t \in [0, T]: \|\Lambda_{\varepsilon}(z_1(\cdot, t)) - \Lambda_{\varepsilon}(z_2(\cdot, t))\|_{L^2(\Omega)} \leq C_\varepsilon \int_0^t \|z_1(\cdot, s) - z_2(\cdot, s)\|_{L^2(\Omega)} \, ds$$

$$+ \frac{1}{\varepsilon} \int_0^t \|g_2(J_\varepsilon z_1(\cdot, s) - J_\varepsilon z_2(\cdot, s))\|_{L^2(\Omega)} \, ds,$$

where $C_\varepsilon \overset{\text{def}}{=} \frac{1}{\varepsilon} (\frac{1}{2} + \|\varepsilon\|_{L^\infty(\Omega)} + C_{zz}^H + C_{zz}^T \|\theta\|_{L^\infty(\Omega, T)})$. Thus the energy estimate (4.13) allows us to infer that

$$\forall s \in [0, T] \forall z \in L^2(0, T; L^2(\Omega)): \|g_2(z(\cdot, s))\|_{L^2(\Omega)} \leq \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \left( \int_0^s \|\text{div}(\varepsilon z(\cdot, \sigma))\|_{L^1(\Omega)} \, d\sigma \right)^{\frac{1}{2}}$$

$$\leq \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \left( \int_0^s \|z(\cdot, \sigma)\|_{L^2(\Omega)}^2 \, d\sigma \right)^{\frac{1}{2}},$$

which, leads to

$$\forall t \in [0, T]: \int_0^t \|g_2(J_\varepsilon z_1(\cdot, s) - J_\varepsilon z_2(\cdot, s))\|_{L^2(\Omega)} \, ds$$

$$\leq \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \int_0^t \left( \int_0^s \|J_\varepsilon z_1(\cdot, \sigma) - J_\varepsilon z_2(\cdot, \sigma)\|_{L^2(\Omega)}^2 \, d\sigma \right)^{\frac{1}{2}} \, ds.$$

Since $J_\varepsilon$ is a contraction on $L^2(\Omega)$ we infer

$$\forall t \in [0, T]: \int_0^t \|g_2(J_\varepsilon z_1(\cdot, s) - J_\varepsilon z_2(\cdot, s))\|_{L^2(\Omega)} \, ds \leq \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \int_0^t \left( \int_0^s \|z_1(\cdot, \sigma) - z_2(\cdot, \sigma)\|_{L^2(\Omega)}^2 \, d\sigma \right)^{\frac{1}{2}} \, ds$$

$$\leq \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} t \|z_1 - z_2\|_{C^0(0, T; L^2(\Omega))}.$$ 

Finally, we find

$$\forall t \in [0, T]: \|\Lambda_{\varepsilon}(z_1(\cdot, t)) - \Lambda_{\varepsilon}(z_2(\cdot, t))\|_{L^2(\Omega)} \leq \left( C_\varepsilon + \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \right) t \|z_1 - z_2\|_{C^0(0, T; L^2(\Omega))}.$$ 

We iterate the previous computation to get

$$\forall t \in [0, T] \forall m \geq 1: \|\Lambda_{\varepsilon}^m(z_1(\cdot, t)) - \Lambda_{\varepsilon}^m(z_2(\cdot, t))\|_{L^2(\Omega)} \leq \left( C_\varepsilon + \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \right)^m \frac{t^m}{m!} \|z_1 - z_2\|_{C^0(0, T; L^2(\Omega))}.$$ 

It follows that there exists $m_0 \geq 1$ such that

$$\left( C_\varepsilon + \frac{\|\varepsilon\|_{L^\infty(\Omega)}^{2}}{\sqrt{c^2 + c_{\text{Korn}}}} \right)^{m_0} \frac{t^{m_0}}{m_0!} < 1,$$

and $\Lambda_{\varepsilon}^{m_0}$ possesses a unique fixed point in $C^0([0, T]; L^2(\Omega))$, which is also the unique fixed point of $\Lambda_{\varepsilon}$. We may conclude that there exists a unique solution $z_\varepsilon \in C^1([0, T]; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}))$ to the problem (4.9)–(4.10).

Let us choose now a sequence $(\theta_\varepsilon)_{\varepsilon > 0}$ such that

$$\theta_\varepsilon \to \theta \text{ in } L^q(0, T; L^p(\Omega)).$$

Since $p, q \geq 2$, it follows that

$$\theta_\varepsilon \to \theta \text{ in } L^2(0, T; L^2(\Omega)).$$

We will establish that there exists a subsequence of $(z_\varepsilon)_{\varepsilon > 0}$, still denoted by $(z_\varepsilon)_{\varepsilon > 0}$, such that

$$z_\varepsilon \rightharpoonup z \text{ in } H^1(0, T; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \text{ weak},$$

$$z_\varepsilon \to z \text{ in } L^\infty(0, T; H^1(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \text{ weak *},$$

where $z$ is a solution of the problem (4.9)–(4.10). To do so, we notice that

$$\forall t \in [0, T]: \dot{z}_\varepsilon(\cdot, t) = (\partial\Psi_{\text{dev}} + M)^{-1} \left( \text{Proj}_{\mathbb{R}^{3 \times 3}_{\text{dev}}} (g(t, J_\varepsilon z_\varepsilon(\cdot, t)) - \partial \varphi_{\varepsilon}(z_\varepsilon(\cdot, t)) - E z_\varepsilon(\cdot, t) - D_z H_1(z_\varepsilon(\cdot, t))) \right),$$

where $D_z H_1$ denotes the derivative of $H_1$ with respect to $z$.
Global existence result for phase transformations with heat transfer in shape memory alloys

Thus, the global existence result for phase transformations with heat transfer in shape memory alloys can be stated as follows:

\[ \forall t \in [0, T] : w_s(t, \Theta, \mathcal{J}_s(z_s(t))) \equiv -(g_1(\Theta_s(\cdot), t) + g_2(\mathcal{J}_s(z_s(t))) + \varphi_0(\cdot, t)D_zH_2(\mathcal{J}_s(z_s(t)))) . \]

Define

\[ \forall t \in [0, T] : \quad w_s(t, \Theta, \mathcal{J}_s(z_s(t))) = \partial \varphi_0(\cdot, t) - \nabla g_s(t, \mathcal{J}_s(z_s(t))) - \mathcal{E}z_s(t) - D_zH_1(z_s(t)) . \]

We have \( w_s \in C^0([0, T]; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{sym}})) \) and

\[ \forall t \in [0, T] : \quad w_s(t, \cdot) = (\mathcal{L}g_s + \mathcal{M})^{-1} \left[ \text{Proj}_{\mathbb{R}^{3 \times 3}}(w_s(t, \cdot)) \right] . \]

In order to obtain a priori estimates, we multiply (4.14a) by \( \mathcal{E}z_s \), and we integrate this expression over \( Q_r \) to get

\[ \int_{Q_r} w_s \dot{z}_s \, dx \, dt + \int_{Q_r} \partial \varphi_0(\cdot, t) \dot{z}_s \, dx \, dt + \int_{Q_r} \mathcal{E}z_s \dot{z}_s \, dx \, dt + \int_{Q_r} D_zH_1(z_s) \dot{z}_s \, dx \, dt = \int_{Q_r} g_s(t, \mathcal{J}_s(z_s)) \dot{z}_s \, dx \, dt . \]

We observe that

\[ \int_{Q_r} \partial \varphi_0(\cdot, t) \dot{z}_s \, dx \, dt = \varphi_0(z_s(\cdot, \tau)) - \varphi_0(z_s^0), \quad \int_{Q_r} D_zH_1(z_s) \dot{z}_s \, dx \, dt = \int_{Q_r} H_1(z_s(\cdot, \tau)) \, dx - \int_{Q_r} H_1(z_s^0) \, dx . \]

Moreover, recalling that \( \dot{z}_s \in C^0([0, T]; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \), we have

\[ \int_{Q_r} w_s \dot{z}_s \, dx \, dt = \int_{Q_r} \text{Proj}_{\mathbb{R}^{3 \times 3}}(w_s) \dot{z}_s \, dx \, dt \geq c_\mathcal{M} \int_0^T \| \dot{z}_s \|^2_{L^2(\Omega)} \, dt . \]

Since \( z_s^0 \in H^1(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}) = \mathcal{D}(\varphi) \), we have \( \varphi_0(z_s^0) \leq \varphi(0) \) (for technical details, the reader is referred to \cite{Bre73}) and

\[ \frac{c_\mathcal{M}}{2} \int_0^T \| \dot{z}_s \|^2_{L^2(\Omega)} \, dt + \varphi_0(z_s(\cdot, \tau)) + \int_\Omega H_1(z_s(\cdot, \tau)) \, dx + \frac{1}{2} \int_\Omega \mathcal{E}z_s(\cdot, \tau) : z_s(\cdot, \tau) \, dx \leq \varphi(0) + \int_\Omega H_1(z_0^0) \, dx + \frac{1}{2} \int_\Omega \mathcal{E}z_0^0 : z_0^0 \, dx + \frac{1}{2c_\mathcal{M}} \int_0^T \| g_s(t, \mathcal{J}_s(z_s(\cdot, \tau))) \|^2_{L^2(\Omega)} \, dt . \]

Furthermore

\[ \forall \epsilon > 0 \forall \tilde{z} \in L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}) : \quad \varphi(\tilde{z}) = \frac{\epsilon}{2} \| \partial \varphi_0(\tilde{z}) \|^2_{L^2(\Omega)} + \varphi(\mathcal{J}_s \tilde{z}) + \frac{\epsilon}{2} \| -z_s \overset{\circ}{-} \mathcal{J}_s \tilde{z} \|^2_{L^2(\Omega)} + \varphi(\mathcal{J}_s \tilde{z}) . \]

Thus \( \mathcal{J}_s \tilde{z} \in C^0([0, T]; H^1(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \) for all \( \tilde{z} \in C^0([0, T]; L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}})) \) and

\[ \frac{c_\mathcal{M}}{2} \int_0^T \| \dot{z}_s \|^2_{L^2(\Omega)} \, dt + \frac{1}{2c_\mathcal{M}} \| z_s(\cdot, \tau) - \mathcal{J}_s z_s(\cdot, \tau) \|^2_{L^2(\Omega)} + \varphi(\mathcal{J}_s z_s(\cdot, \tau)) + \int_\Omega H_1(z_s(\cdot, \tau)) \, dx \leq \varphi(0) + \int_\Omega H_1(z_0^0) \, dx + \frac{1}{2} \int_\Omega \mathcal{E}z_0^0 : z_0^0 \, dx + \frac{1}{2c_\mathcal{M}} \int_0^T \| g_s(t, \mathcal{J}_s z_s(\cdot, \tau)) \|^2_{L^2(\Omega)} \, dt . \]

By using (3.11a), (3.12) and (4.13a), we infer that

\[ \frac{c_\mathcal{M}}{2} \int_0^T \| \dot{z}_s \|^2_{L^2(\Omega)} \, dt + C_1 \| \mathcal{J}_s z_s(\cdot, \tau) \|^2_{L^2(\Omega)} + \| z_s(\cdot, \tau) \|^2_{L^2(\Omega)} \leq C_0 + \frac{1}{2c_\mathcal{M}} \int_0^T \| g_s(t, \mathcal{J}_s z_s(\cdot, \tau)) \|^2_{L^2(\Omega)} \, dt , \]

where \( C_0 \equiv \varphi(0) + \int_\Omega H_1(z_0^0) \, dx + \frac{1}{2} \int_\Omega \mathcal{E}z_0^0 : z_0^0 \, dx \) and \( C_1 \equiv \min(c_{H^1} + \frac{\epsilon}{2}, \frac{\epsilon}{2}) \). Since \( \mathcal{J}_s \) is a contraction and \( \mathcal{J}_s 0 = 0 \), we have

\[ \forall \tilde{z} \in L^2(\Omega; \mathbb{R}^{3 \times 3}_{\text{dev}}) : \quad \| \mathcal{J}_s \tilde{z} \|^2_{L^2(\Omega)} = \| \mathcal{J}_s \tilde{z} - \mathcal{J}_s 0 \|^2_{L^2(\Omega)} \leq \| \tilde{z} \|^2_{L^2(\Omega)} . \]

Then we find

\[ \frac{c_\mathcal{M}}{2} \int_0^T \| \dot{z}_s \|^2_{L^2(\Omega)} \, dt + C_1 \| \mathcal{J}_s z_s(\cdot, \tau) \|^2_{H^1(\Omega)} \leq C_0 + \frac{1}{2c_\mathcal{M}} \int_0^T \| g_s(t, \mathcal{J}_s z_s(\cdot, \tau)) \|^2_{L^2(\Omega)} \, dt . \]
For the bound on $g_\epsilon(t, \mathcal{J}_\epsilon z(\cdot, t))$, we use (3.12) to get

$$
\int_0^T \|g_\epsilon(t, \mathcal{J}_\epsilon z)\|_{L^2(\Omega)}^2 \, dt \leq 3 \int_0^T \|g_1(\theta_\epsilon)\|_{L^2(\Omega)}^2 \, dt + 3 \int_0^T \|g_2(\mathcal{J}_\epsilon z)\|_{L^2(\Omega)}^2 \, dt 
+ 6(C_2^H)^2 \int_0^T \|\theta_\epsilon\|_{L^2(\Omega)}^2 \, dt + 6(C_2^H)^2 \int \mathcal{J}_\epsilon z \, dz \, dt.
$$

By using the definition of mappings $g_1$ and $g_2$, the first two terms on the right hand side of (4.17) can be estimated. More precisely, it is plain that there exists $C_2 > 0$ such that

$$
\int_0^T \|g_1(\theta_\epsilon)\|_{L^2(\Omega)}^2 \, dt + \int_0^T \|g_2(\mathcal{J}_\epsilon z)\|_{L^2(\Omega)}^2 \, dt
\leq C_2 \left( \|u_0\|_{H^1(\Omega)}^2 + \int_0^\infty \|\theta_\epsilon\|_{L^2(\Omega)}^2 \, dt + \int_0^\infty \|\theta_\epsilon\|_{L^2(\Omega)}^2 \, dt + \int_0^\infty \|\mathcal{J}_\epsilon z\|_{L^2(\Omega)}^2 \, dt \right).
$$

The last term on the right hand side of (4.17) is estimated by using Hölder’s inequality and the continuous embedding $H^1(\Omega) \hookrightarrow L^4(\Omega)$. Namely, it follows that there exists $C_3 > 0$ such that

$$
\int_{Q^*} \|\theta_\epsilon\|_{L^2(\Omega)}^2 |\mathcal{J}_\epsilon z| \, d\tau \, dt \leq C_3 \int_0^\infty \|\theta_\epsilon\|_{L^2(\Omega)}^2 \|\mathcal{J}_\epsilon z\|_{H^1(\Omega)}^2 \, dt.
$$

We insert (4.18) and (4.19) into (4.16), we obtain

$$
\frac{1}{2} \int_0^T \|\tilde{z}_\epsilon\|_{L^2(\Omega)}^2 \, dt + C_1 \|\mathcal{J}_\epsilon z(\cdot, \tau)|_{H^1(\Omega)}^2 \leq C_0 + \frac{C_2^2}{2\epsilon} \left( \|u_0\|_{H^1(\Omega)}^2 + \int_0^T \|\ell\|_{L^2(\Omega)}^2 \, dt \right) 
+ \frac{C_2^2 + 6(C_2^H)^2}{2\epsilon} \int_0^T \|\theta_\epsilon\|_{L^2(\Omega)}^2 \, dt + \int_0^T \|\mathcal{J}_\epsilon z\|_{H^1(\Omega)}^2 \, dt.
$$

Since $\theta_\epsilon$ is bounded in $L^q(0, T; L^p(\Omega))$, with $q \geq 2$ and $p \geq 4$, we may deduce from Grönwall’s lemma that $\mathcal{J}_\epsilon z$ is bounded in $L^\infty(0, T; H^1(\Omega))$, $g_\epsilon(\cdot, \mathcal{J}_\epsilon z)$ is bounded in $L^2(0, T; L^6(\Omega))$ and $z_\epsilon$ is bounded in $H^1(0, T; L^2(\Omega)) \cap L^\infty(0, T; L^2(\Omega))$, independently of $\epsilon > 0$. Therefore (3.12) implies that $D_z H_1(z_\epsilon)$ is bounded in $L^\infty(0, T; L^2(\Omega))$, independently of $\epsilon > 0$. Furthermore the definition of the subdifferential $\partial \Psi_{dev}(z_\epsilon)$ enables us to infer from (4.14b) that

$$
- \int_{Q^*} M \tilde{z}_\epsilon w^p \, dx \, dt + \int_0^T \|w^p\|_{L^2(\Omega)}^2 \, dt \leq \int_{Q^*} (\Psi(w^p + \tilde{z}_\epsilon) - \Psi(\tilde{z}_\epsilon)) \, dx \, dt \leq \int_{Q^*} \Psi(w^p) \, dx \, dt,
$$

where $w^p = \text{Proj}_{H^1(\Omega)}(w_\epsilon)$, and using (3.10b), we get

$$
\|w^p\|_{L^2(0, T; L^2(\Omega))} \leq C^\Psi \sqrt{T|\Omega|} + C^M \|\tilde{z}_\epsilon\|_{L^2(0, T; L^2(\Omega))}.
$$

Observe that $\tilde{z}_\epsilon$ is bounded in $L^2(0, T; L^2(\Omega))$, independently of $\epsilon > 0$, allows us to conclude that $w^p$ is bounded in $L^2(0, T; L^2(\Omega))$, independently of $\epsilon > 0$. On the other hand, we multiply (4.14b) by $\partial \varphi_\epsilon(z_\epsilon)$ and we integrate this result over $Q^*$. Recalling that $\partial \varphi_\epsilon(z_\epsilon)$ takes its values in $L^2(\Omega; \mathbb{R}^{3 \times 3})$, the Cauchy-Schwarz’s inequality, (3.12) and (4.21) give

$$
\|\partial \varphi_\epsilon(z_\epsilon)\|_{L^2(0, T; L^2(\Omega))} \leq (C_H^1 + C^\Psi \sqrt{T|\Omega|}) + (C_H^1 + \|\xi\|_{L^\infty(\Omega)}) \|z_\epsilon\|_{L^2(0, T; L^2(\Omega))} 
+ C^M \|\tilde{z}_\epsilon\|_{L^2(0, T; L^2(\Omega))} + \|g_\epsilon(\cdot, \mathcal{J}_\epsilon z)\|_{L^2(0, T; L^2(\Omega))}.
$$

Thus $\partial \varphi_\epsilon(z_\epsilon)$ is bounded in $L^2(0, T; L^2(\Omega))$, independently of $\epsilon > 0$ and finally $w_\epsilon$ is bounded in $L^2(0, T; L^2(\Omega))$, independently of $\epsilon > 0$. Furthermore, there exists $z \in H^1(0, T; L^2(\Omega)) \cap L^\infty(0, T; L^2(\Omega)) \cap L^\infty(0, T; H^1(\Omega)), \tilde{z} \in L^\infty(0, T; H^1(\Omega)), w, v \in L^2(0, T; L^2(\Omega))$ such that it is possible to extract subsequences, still denoted by $z_\epsilon$, $\mathcal{J}_\epsilon z_\epsilon$, $w_\epsilon$ and $\partial \varphi_\epsilon(z_\epsilon)$ satisfying the following convergences

$$
\begin{align*}
z_\epsilon & \to z \quad \text{in} \quad H^1(0, T; L^2(\Omega)) \quad \text{weak}, \\
\mathcal{J}_\epsilon z_\epsilon & \to \tilde{z} \quad \text{in} \quad L^\infty(0, T; L^2(\Omega)) \quad \text{weak *}, \\
w_\epsilon & \to w \quad \text{in} \quad L^2(0, T; L^2(\Omega)) \quad \text{weak}, \\
\partial \varphi_\epsilon(z_\epsilon) & \to v \quad \text{in} \quad L^2(0, T; L^2(\Omega)) \quad \text{weak}.
\end{align*}
$$
Moreover, using (4.13) there exists $C_4 > 0$ such that
\[
\forall \tau \in [0, T]: \|z_\epsilon(\cdot, \tau) - J_\epsilon z_\epsilon(\cdot, \tau)\|_{L^2(\Omega)}^2 \leq C_4 \epsilon (1 + \|g_\epsilon(\cdot, J_\epsilon z_\epsilon)\|_{L^2(0, T; L^2(\Omega))}).
\] (4.22)
which, allows us to infer that $z = \bar{z}$. Since $J_\epsilon$ is a contraction on $L^2(\Omega)$, it follows also that, possibly extracting another subsequence, still denoted by $J_\epsilon z_\epsilon$, we have
\[
\forall \tau \in [2, 6]: J_\epsilon z_\epsilon \to z \text{ in } C^0([0, T]; L^2(\Omega)).
\]
We may deduce from (4.22) that
\[
z_\epsilon \to z \text{ in } C^0(0, T; L^2(\Omega)).
\]
Our aim now consists to pass to the limit in (4.14). To do so, we observe that the mapping $L_0$ is linear and continuous, the mappings $g_1$ and $g_2$ are also continuous, which, gives
\[
g_1(\theta_\epsilon) \to g_1(\theta) \text{ in } L^2(0, T; L^2(\Omega)),
g_2(J_\epsilon z_\epsilon) \to g_2(z) \text{ in } L^2(0, T; L^2(\Omega)).
\]
Since $D_z H_i$, $i = 1, 2$, is Lipschitz continuous, it follows that
\[
D_z H_1(z_\epsilon) + E z_\epsilon \to D_z H_1(z) + E z \text{ in } C^0([0, T]; L^2(\Omega)),
\]
\[
\forall \tau \in [2, 6]: D_z H_2(J_\epsilon z_\epsilon) \to D_z H_2(z) \text{ in } C^0([0, T]; L^2(\Omega)).
\]
Hence we consider $r$ such that $\frac{1}{p} + \frac{1}{r} = \frac{1}{2}$ to get
\[
\theta_d D_z H_2(J_\epsilon z_\epsilon) \to \theta_d D_z H_2(z) \text{ in } L^2(0, T; L^2(\Omega)).
\]
We can pass to the limit in all terms of (4.14a), we obtain
\[
w + v + D_z H_1(z) + E z = -(g_1(\theta) + g_2(z) + \theta_d D_z H_2(z)).
\] (4.23)
It remains to pass to the limit in (4.14b) and to prove that $z$ solves (4.5)–(4.7). Using (4.10) and the strong convergence of $(z_\epsilon)_{\epsilon > 0}$ to $z$ in $C^0([0, T]; L^2(\Omega))$, we infer that $z(0, \cdot) = z^0$. Then we prove that $v(t, \cdot) \in \partial \varphi(z(t, \cdot))$ for almost every $t \in [0, T]$, i.e., $v \in \partial \varphi(z)$ where $\partial \varphi$ is identified to its $L^2(0, T)$-extension. To do so, we use the classical results for maximal monotone operators (see [Bre73]). More precisely, since $\partial \varphi(z_\epsilon) \in \partial \varphi(J_\epsilon z_\epsilon)$, it is sufficient to prove that
\[
\limsup_{\epsilon \to 0} \int_{Q_T} (\partial \varphi(z_\epsilon)) : (J_\epsilon z_\epsilon) \, dx \, dt \leq \int_{Q_T} w : z \, dx \, dt,
\] (4.24)
which is an immediate consequence of the convergence results for the subsequences $(\partial \varphi(z_\epsilon))_{\epsilon > 0}$ and $(J_\epsilon z_\epsilon)_{\epsilon > 0}$. Hence, using the definitions of $w$ and $\varphi$, we may conclude that $z$ is a solution of (4.5)–(4.7) if we can prove that $\text{Proj}_{\mathcal{J}_{\text{dev}}^3}^{3 \times 3}(w(t, \cdot)) - \Psi_{\text{dev}}(\dot{z}(t, \cdot))$ for almost every $t \in [0, T]$. So we may use the same trick as previously and we only need to check that
\[
\limsup_{\epsilon \to 0} \int_{Q_T} \text{Proj}_{\mathcal{J}_{\text{dev}}^3}^{3 \times 3}(w_\epsilon) : \dot{z}_\epsilon \, dx \, dt \leq \int_{Q_T} \text{Proj}_{\mathcal{J}_{\text{dev}}^3}^{3 \times 3}(w) : \dot{z} \, dx \, dt.
\] (4.25)
But $\dot{z}_\epsilon$ takes its values in $\mathbb{R}^{3 \times 3}_{\text{dev}}$ and thus $\dot{z}$ takes also its values in $\mathbb{R}^{3 \times 3}_{\text{dev}}$. It follows that (4.25) is equivalent to
\[
\limsup_{\epsilon \to 0} \int_{Q_T} w_\epsilon : \dot{z}_\epsilon \, dx \, dt \leq \int_{Q_T} w : \dot{z} \, dx \, dt.
\]
We compute $\int_{Q_T} w_\epsilon : \dot{z}_\epsilon \, dx \, dt$ and $\int_{Q_T} w : \dot{z} \, dx \, dt$ from (4.14a) and (4.23), then the convergence results obtained above imply that (4.25) holds if and only if
\[
\liminf_{\epsilon \to 0} \int_{Q_T} \partial \varphi(z_\epsilon) : \dot{z}_\epsilon \, dx \, dt \geq \int_{Q_T} \partial \varphi(z) : \dot{z} \, dx \, dt.
\]
We observe that \( \int_{\Omega_T} \partial \varphi(z_t) \cdot z_t \, dx \, dt = \varphi(z_t(T)) - \varphi(z_0) \geq \varphi(J_z z_t(T)) - \varphi(z_0) \) and \( \int_{\Omega_T} \partial \varphi(z_t) \cdot z_t \, dx \, dt = \varphi(z(T)) - \varphi(z_0) \). Recalling that \( z_0 \in D(\varphi) \), we get \( \lim_{t \to 0} \varphi(z_t) = \varphi(z_0) \) and the lower semicontinuity of \( \varphi \), allows us to conclude. This proves the existence result.

Finally we observe that
\[
-\text{Proj}_{\text{dev}}^{3 \times 3}(w) = \text{Proj}_{\text{dev}}^{3 \times 3}(g_1(\theta) + g_2(z) + \mathbb{E} z) + D_z H_1(z) + \theta D_z H_2(z) + \nu, \\
\]
and, using the definition of the mappings \( g_1, g_2 \) and \( \varphi \)
\[
-\text{Proj}_{\text{dev}}^{3 \times 3}(w) = \text{Proj}_{\text{dev}}^{3 \times 3}(-\mathbb{E}(e(u) - z)) + D_z H_1(z) + \theta D_z H_2(z) - \nu \Delta z.
\]
So we may rewrite (4.11) as follows
\[
\dot{z} - \nu M^{-1} \Delta z = M^{-1} f^z,
\]
with \( f^z \equiv \text{Proj}_{\text{dev}}^{3 \times 3}(\mathbb{E}(e(u) - z)) - D_z H_1(z) - \theta D_z H_2(z) - \psi \) and \( \psi = \text{Proj}_{\text{dev}}^{3 \times 3}(w) - M \dot{z} \in \partial \Psi_{\text{dev}}(\dot{z}). \)

With (4.10) we infer that
\[
\forall \psi \in \partial \Psi_{\text{dev}}(\dot{z}) : |\psi| \leq C^\Psi \text{ a.e. } (x, t) \in \Omega \times (0, T).
\]
Furthermore, since \( z \in L^\infty(0, T; H^1(\Omega)) \), we infer with (3.12) that \( D_z H_i(z) \) belongs to \( L^\infty(0, T; L^q(\Omega)) \) for \( i = 1, 2 \). Then it follows that \( M^{-1} f^z \) belongs to \( L^3(0, T; L^q(\Omega)) \). Since \( z_0 \in X_{q,p}(\Omega) \), we may deduce from the maximal regularity result for parabolic systems that \( z \in L^3(0, T; H^1(\Omega)) \) and \( z \in L^2(0, T; L^2(\Omega)) \). We refer to [Dor93, HR98, PS01] and the references therein for more details than the maximal regularity result for parabolic systems and its consequences.

Let us observe that here neither \( \partial \Psi_{\text{dev}} + M \) nor \( \partial \varphi + D_z H_1 + \theta D_z H_2 + \text{Proj}_{\text{dev}}^{3 \times 3}(\mathbb{E} + g_1 + g_2) \) are linear and self-adjoint operators of \( L^2(\Omega; \mathbb{R}^{3 \times 3}) \), so we cannot use the ideas proposed in [CoV90] to prove uniqueness. The uniqueness result proved below relies on the boundedness assumption (4.11) for the hardening functional \( H_1 \) combined with Grönwall’s lemma. More precisely let \( h_1 \in C^2(\mathbb{R}_{\text{dev}}; \mathbb{R}) \) be defined by
\[
\forall z \in \mathbb{R}_{\text{dev}}^{3 \times 3} : h_1(z) \overset{\text{def}}{=} H_1(z) - C^{H_1}|z|^2,
\]
with a real number \( C^{H_1} > 0 \). Assumption (4.11) implies that there exists \( C^{h_1} > 0 \) such that
\[
\forall z_1, z_2 \in \mathbb{R}_{\text{dev}}^{3 \times 3} : |D_z h_1(z_1) - D_z h_1(z_2)| \leq C^{h_1}|z_1 - z_2|.
\]

**Proposition 4.2 (Uniqueness for \((P_{u,T})\))** Assume that \( \theta \) is given in \( L^q(0, T; L^p(\Omega)) \), (3.10), (3.11), (3.14) and (3.16) hold and \( u^0 \in H^1(\Omega) \), \( z^0 \in X_{q,p}(\Omega) \). Then the problem (4.1)–(4.3) admits a unique solution.

**Proof.** Let \( \xi_1 \overset{\text{def}}{=} (u_1, z_1) \) and \( \xi_2 \overset{\text{def}}{=} (u_2, z_2) \) be two solutions of (4.1a)–(4.1b) satisfying (4.2) and (4.3). With the results of the previous theorem we already know that \( (u_i, z_i) \in C^0([0, T]; H_0^1(\Omega) \times H^1(\Omega)) \) and \( \Delta z_i \in L^3(0, T; L^2(\Omega)) \), \( i = 1, 2 \). Define
\[
\gamma(t) \overset{\text{def}}{=} \frac{1}{2} \int_\Omega \mathbb{E}((e(u_1) - z_1) - (e(u_2) - z_2))(e(u_1) - z_1) - (e(u_2) - z_2)dx \\
+ \frac{\nu}{2} \int_\Omega \nabla (z_1 - z_2) : \nabla (z_1 - z_2) dx + C^{H_1} \int_\Omega |z_1 - z_2|^2 dx
\]
for all \( t \in [0, T] \). Since
\[
\gamma(t) = \frac{1}{2} \int_\Omega \mathbb{E}((e(u_1) - z_1) - (e(u_2) - z_2))(e(u_1) - z_1) - (e(u_2) - z_2)dx \\
- \frac{\nu}{2} \int_\Omega \Delta (z_1 - z_2) : (z_1 - z_2) dx + C^{H_1} \int_\Omega |z_1 - z_2|^2 dx,
\]
the mapping \( \gamma \) is continuous on \([0, T]\) and its derivative in the sense of distributions belongs to \( L^1(0, T) \). Thus \( \gamma \) is absolutely continuous on \([0, T]\) and with assumptions (4.12) and (4.13) combined with Korn’s inequality, we infer that there exists a real number \( \kappa > 0 \) such that
\[
\gamma(t) \geq \kappa (\|u_1(\cdot, t) - u_2(\cdot, t)\|_{H^1(\Omega)}^2 + \|z_1(\cdot, t) - z_2(\cdot, t)\|_{H^1(\Omega)}^2) \forall t \in [0, T].
\]
On the one hand, recalling that $\partial \Psi_{\text{dev}}(\cdot)$ is a monotone operator, the Green’s formula and (4.3) enable us to deduce from (4.11) that

$$\forall \tilde{z}_i \in L^2(\Omega; \mathbb{R}^{3\times 3}): \int_{\Omega} (M\tilde{z}_i - \mathcal{E}(e(\bar{u}_i) - \bar{z}_i)) + D_xH_1(z_1) + \theta D_xH_2(z_1)(\tilde{z}_i - \bar{z}_i) \, dx$$

$$- \nu \int_{\Omega} \Delta \tilde{z}_i(\tilde{z}_i - \bar{z}_i) \, dx \leq \int_{\Omega} (\Psi_{\text{dev}}(\tilde{z}_i) - \Psi_{\text{dev}}(\bar{z}_i)) \, dx. \tag{4.31}$$

On the other hand, we multiply (4.11) by $\dot{u}_i - \bar{u}_i$, we integrate this expression over $\Omega$, and with the help of the Green’s formula together with (4.3), we obtain

$$\forall \bar{u}_i \in H^1_0(\Omega): \int_{\Omega} (\mathcal{E}(e(\bar{u}_i) - \bar{z}_i) + \alpha \theta I + \mathbb{L} e(\bar{u}_i))(e(\bar{u}_i) - e(\bar{u}_i)) \, dx = \int_{\Omega} \ell(\dot{u}_i - \bar{u}_i) \, dx. \tag{4.32}$$

Thus we add (4.31) and (4.32), we get

$$\int_{\Omega} (M\tilde{z}_i - \mathcal{E}(e(\bar{u}_i) - \bar{z}_i)) + D_xH_1(z_1) + \theta D_xH_2(z_1)(\tilde{z}_i - \bar{z}_i) \, dx - \nu \int_{\Omega} \Delta \tilde{z}_i(\tilde{z}_i - \bar{z}_i) \, dx$$

$$+ \int_{\Omega} (\mathcal{E}(e(\bar{u}_i) - \bar{z}_i) + \alpha \theta I + \mathbb{L} e(\bar{u}_i))(e(\bar{u}_i) - e(\bar{u}_i)) \, dx \leq \int_{\Omega} (\ell(\dot{u}_i - \bar{u}_i) + \Psi_{\text{dev}}(\tilde{z}_i) - \Psi_{\text{dev}}(\bar{z}_i)) \, dx. \tag{4.33}$$

We choose now $(\bar{u}_i, \tilde{z}_i) = (\bar{u}_{3-i}, \tilde{z}_{3-i})$ for $i = 1, 2$ in (4.33), we add these two inequalities, and with the help of (3.14), we find

$$\dot{\gamma}(t) + c^M \int_{\Omega} |\dot{z}_1 - \dot{z}_2|^2 \, dx + c^L \int_{\Omega} |e(\bar{u}_1) - e(\bar{u}_2)|^2 \, dx \leq - \int_{\Omega} (D_xh_1(z_1) - D_xh_1(z_2))(\dot{z}_1 - \dot{z}_2) \, dx - \int_{\Omega} \theta (D_xH_2(z_1) - D_xH_2(z_2))(\dot{z}_1 - \dot{z}_2) \, dx \quad \text{a.e. } t \in [0, T].$$

We estimate the terms of the right hand side of the previous inequality as follows:

$$\left| \int_{\Omega} (D_xh_1(z_1) - D_xh_1(z_2))(\dot{z}_1 - \dot{z}_2) \, dx \right| \leq C_{h_1} \int_{\Omega} |z_1 - z_2| |\dot{z}_1 - \dot{z}_2| \, dx$$

$$\leq \frac{M}{4} \int_{\Omega} |\dot{z}_1 - \dot{z}_2|^2 \, dx + \frac{(C_{h_1})^2}{c^L} \int_{\Omega} |z_1 - z_2|^2 \, dx,$$

and

$$\left| \int_{\Omega} \theta (D_xH_2(z_1) - D_xH_2(z_2))(\dot{z}_1 - \dot{z}_2) \, dx \right| \leq C_{H_2} \int_{\Omega} |\theta| |z_1 - z_2| |\dot{z}_1 - \dot{z}_2| \, dx$$

$$\leq \frac{M}{4} \int_{\Omega} |\dot{z}_1 - \dot{z}_2|^2 \, dx + \frac{(C_{H_2})^2}{c^L} \int_{\Omega} |\theta|^2 |z_1 - z_2|^2 \, dx + \frac{C_{H_2}^2}{c^L} \|\theta\|_{L^4(\Omega)}^2 \|z_1 - z_2\|^2_{L^4(\Omega)}.$$
Proof. We consider \( \vartheta_i \in L^q(0, T; L^p(\Omega)) \) and for \( i = 1, 2 \), we denote by \( \vartheta_i \overset{\text{def}}{=} \zeta(\vartheta_i) \in L^q(0, T; L^p(\Omega)) \) and \((u_i, z_i)\) the solution of the following system:

\[
\begin{align*}
- \text{div}(\mathcal{E}(e(u_i)) - z_i) + \alpha \vartheta_i \text{I} + \text{Le}(u_i) &= \ell, \\
\partial \Psi(z_i) + M \dot{z}_i - \mathcal{E}(e(u_i)) = D_z H_1(z_i) + \vartheta D_z H_2(z_i) - \nu \Delta z_i &\ni 0,
\end{align*}
\]

(4.34a) (4.34b)

together with initial conditions

\[
u(0, \cdot) = u^0, \quad \dot{z}(0, \cdot) = z^0,
\]

(4.35)

and boundary conditions

\[
u_{|\bar{\Omega}} = 0, \quad \nabla z_{|\bar{\Omega}} = 0.
\]

(4.36)

Since the mapping \( \vartheta \mapsto \theta = \zeta(\vartheta) \) is continuous from \( L^q(0, T; L^p(\Omega)) \) to \( L^q(0, T; L^p(\Omega)) \), it is enough to check that the mapping \( \theta = \zeta(\vartheta) \mapsto (u, z) \) is continuous from \( L^q(0, T; L^p(\Omega)) \) to \( L^q(0, T; L^p(\Omega)) \).

Once again the key tool is the structural decomposition (4.27) combined with Grönwall’s lemma. We reproduce the same kind of computations as in Proposition 4.2. More precisely, the total energy inequality associated to (4.34) is obtained by multiplying (4.34a) and (4.34b) by \( \dot{u}_{3-i} - \dot{u}_i \) and \( \dot{z}_{3-i} - \dot{z}_i \) respectively, and integrating over \( \Omega \). Then, we add these two inequalities, we find

\[
\int_{\Omega} \mathcal{E}(e(u_i)) - z_i + \theta D_z H_2(z_i) - \nu \Delta z_i \, dx + \int_{\Omega} M \dot{z}_i + (\dot{z}_{3-i} - \dot{z}_i) \, dx 
\]

\[
+ \int_{\Omega} \mathcal{E}(e(u_i)) - z_i + \theta D_z H_2(z_i) - \nu \Delta z_i \, dx + \int_{\Omega} M \dot{z}_i + (\dot{z}_{3-i} - \dot{z}_i) \, dx 
\]

\[
+ \int_{\Omega} \theta D_z H_2(z_i) - \dot{z}_i \, dx - \int_{\Omega} \ell (\dot{u}_{3-i} - \dot{u}_i) + \Psi_{\text{dev}}(\dot{z}_{3-i} - \dot{z}_i) + 0 \text{ a.e. } t \in [0, T].
\]

(4.37)

It is convenient to introduce the notations: \( \ddot{u} \overset{\text{def}}{=} u_{1-2}, \ddot{z} \overset{\text{def}}{=} z_{1-2} \) and \( \ddot{\theta} \overset{\text{def}}{=} \theta_1 - \theta_2 \). Therefore, we take \( i = 1, 2 \) in (4.37) and add these two inequalities, we get

\[
\int_{\Omega} \mathcal{E}(e(\ddot{u}) - \ddot{z}) : (e(\ddot{u}) - \ddot{z}) \, dx + \int_{\Omega} M \ddot{z} : \ddot{z} \, dx + \int_{\Omega} \mathcal{E}(e(\ddot{u})) : (e(\ddot{u}) - \ddot{z}) \, dx - \int_{\Omega} \nu \Delta \ddot{z} : \ddot{z} \, dx 
\]

\[
+ \int_{\Omega} (D_z H_1(z_1) - D_z h_1(z_2)) : \ddot{z} \, dx \leq - \int_{\Omega} \alpha \dot{\theta} tr(e(\ddot{u})) \, dx - \int_{\Omega} (\theta_1 D_z H_2(z_1) - \theta_2 D_z H_2(z_2)) : \ddot{z} \, dx.
\]

Define

\[
\forall t \in [0, T] : \gamma(t) \overset{\text{def}}{=} \frac{1}{2} \int_{\Omega} \mathcal{E}(e(u_1) - z_1) - (e(u_2) - z_2) : ((e(u_1) - z_1) - (e(u_2) - z_2)) \, dx 
\]

\[
- \frac{\nu}{2} \int_{\Omega} \nabla(z_1 - z_2) : \nabla(z_1 - z_2) + C^{H_1} \int_{\Omega} |z_1 - z_2|^2 \, dx.
\]

As in Proposition 4.2 we can check that \( \gamma \) is absolutely continuous on \([0, T]\) and (3.14) and (4.27) imply that

\[
\dot{\gamma}(t) + c^M \| \dot{z} \|^2_{L^2(\Omega)} + c^f \| e(\ddot{u}) \|^2_{L^2(\Omega)} \leq - \int_{\Omega} (D_z h_1(z_1) - D_z h_1(z_2)) : \ddot{z} \, dx 
\]

\[
- \int_{\Omega} \alpha \dot{\theta} tr(e(\ddot{u})) \, dx - \int_{\Omega} (\theta_1 D_z H_2(z_1) - \theta_2 D_z H_2(z_2)) : \ddot{z} \, dx,
\]

for almost every \( t \in [0, T] \). Clearly, it follows from (4.28) that

\[
\dot{\gamma}(t) + c^M \| \dot{z} \|^2_{L^2(\Omega)} + c^f \| e(\ddot{u}) \|^2_{L^2(\Omega)} \leq - \int_{\Omega} \alpha \dot{\theta} tr(e(\ddot{u})) \, dx 
\]

\[
- \int_{\Omega} (\theta_1 D_z H_2(z_1) - \theta_2 D_z H_2(z_2)) : \ddot{z} \, dx + C^{H_1} \int_{\Omega} |\dot{z}| \| \dot{z} \| \, dx.
\]

We estimate the first and third term on the right hand side with the help of Cauchy-Schwarz’s inequality, while for the second term, we use the following decomposition

\[
(\theta_1 D_z H_2(z_1) - \theta_2 D_z H_2(z_2)) : \ddot{z} = (\tilde{\theta} D_z H_2(z_1) + \theta_2 (D_z H_2(z_1) - D_z H_2(z_2)) : \ddot{z}.
\]
Hence, we obtain
\[
\dot{\gamma}(t) + \frac{3\gamma}{2}\|\ddot{z}\|_{L^2(\Omega)}^2 + \frac{\dot{\gamma}}{T}\|\sigma\|_{L^2(\Omega)}^2 \leq \left(\frac{C_{h_{1/2}}}{\kappa}\right)^2 \|\dot{\theta}\|_{L^2(\Omega)}^2 + \frac{3\gamma^2}{2T}\|\dot{\theta}\|_{L^2(\Omega)}^2 + \int_\Omega (|\theta||D_zH_2(z_1)||\dot{z}| + |\theta||D_zH_2(z_1) - D_zH_2(z_2)||\dot{z}|) \, dx.
\]
(4.38)

It remains to estimate the last term on the right hand side of (4.38). We use (3.11b) and (3.12) to get
\[
\int_\Omega (|\theta||D_zH_2(z_1)||\dot{z}| + |\theta||D_zH_2(z_1) - D_zH_2(z_2)||\dot{z}|) \, dx \leq C_{H_2}^2 \int_\Omega (1 + |z_1|)\|\dot{z}\| \, dx + C_{zz}^2 \int_\Omega |\theta||\dot{z}| \, dx.
\]
The Young’s inequality implies that there exists \(\gamma_i > 0\), \(i = 1, 2, 3\), such that
\[
\int_\Omega (|\theta||D_zH_2(z_1)||\dot{z}| + |\theta||D_zH_2(z_1) - D_zH_2(z_2)||\dot{z}|) \, dx \leq \frac{C_{H_2}^2}{\gamma_i} \|\theta\|_{L^2(\Omega)}^2 + \frac{C_{zz}^2}{\gamma_3} \int_\Omega \theta^2 \|\dot{z}\|^2 \, dx + \frac{C_{H_2}^2\gamma_2 + C_{H_2}^2\gamma_3}{2} \|\dot{z}\|_{L^2(\Omega)}^2.
\]
(4.39)

Since \(z_1 \in L^2(0, T; H^2(\Omega))\) and \(H^2(\Omega) \hookrightarrow L^\infty(\Omega)\) with continuous embedding, we get
\[
\int_\Omega (|\theta||D_zH_2(z_1)||\dot{z}| + |\theta||D_zH_2(z_1) - D_zH_2(z_2)||\dot{z}|) \, dx \leq \frac{C_{H_2}^2}{\gamma_i} \|\theta\|_{L^2(\Omega)}^2 + \frac{C_{zz}^2}{\gamma_3} \int_\Omega \theta^2 \|\dot{z}\|^2 \, dx + \frac{C_{H_2}^2\gamma_2 + C_{H_2}^2\gamma_3}{2} \|\dot{z}\|_{L^2(\Omega)}^2.
\]
We insert (4.39) in (4.38) and we choose \(\gamma_1 = \gamma_2 = \frac{3\gamma}{4C_{zz}^2}\) and \(\gamma_3 = \frac{3\gamma}{2C_{zz}^2}\). Therefore the continuous embedding \(H^1(\Omega) \hookrightarrow L^4(\Omega)\) and (4.30) give
\[
\dot{\gamma}(t) + \frac{3\gamma}{2}\|\ddot{z}\|_{L^2(\Omega)}^2 + \frac{\dot{\gamma}}{T}\|\sigma\|_{L^2(\Omega)}^2 \leq C(\bar{\theta}, z_1) + C(\theta_2)\|\dot{z}\|_{H^1(\Omega)}^2 \leq C(\bar{\theta}, z_1) + \frac{C(\theta_2)}{\kappa}\gamma(t),
\]
for almost every \(t \in [0, T]\). Here \(C(\bar{\theta}, z_1) \overset{\text{def}}{=} \left(\frac{3\gamma}{2}\right)^2 + \frac{2(C_{H_2})^2}{\kappa}\|\dot{\theta}\|_{L^2(\Omega)}^2 + \frac{2(C_{H_2})^2}{\kappa}\|\theta\|_{H^1(\Omega)}^2 \|\dot{z}\|_{L^2(\Omega)}^2 \|\|z_1\|_{L^\infty(\Omega)}^2 \|\dot{\theta}\|_{L^2(\Omega)}^2\) and \(C(\theta_2) \overset{\text{def}}{=} \left(\frac{5\gamma}{2}\right)^2 + \frac{2(C_{H_2})^2}{\kappa}\|\theta_2\|_{L^4(\Omega)}^2\) where \(C > 0\) is the generic constant involved in the continuous embedding \(H^1(\Omega) \hookrightarrow L^4(\Omega)\). Since \(\bar{q} \geq 3\) we can check that \(C(\bar{\theta}, z_1) \in L^1(0, T)\) and
\[
\int_0^T C(\bar{\theta}, z_1) \, dt \leq \left(\frac{3\gamma}{2}\right)^2 + \frac{2(C_{H_2})^2}{\kappa}\|\dot{\theta}\|_{L^2(0, T; L^2(\Omega))}^2 + \frac{2(C_{H_2})^2}{\kappa}\|\theta_2\|_{L^4(0, T; L^4(\Omega))}^2 \|\|z_1\|_{L^2(0, T; L^2(\Omega))}^2 \|\|\dot{\theta}\|_{L^2(0, T; L^2(\Omega))}^2\).
\]
Note that \(C(\theta_2) \in L^1(0, T)\), which, thanks to Grönwall’s lemma, gives
\[
\forall t \in [0, T]: \quad \gamma(t) \leq \int_0^t C(\bar{\theta}(\cdot, s), z_1(\cdot, s))e^{\frac{\dot{\gamma}}{T} \int_0^s C(\bar{\theta}(\cdot, s) - \bar{\theta}(\cdot, s)) \, ds} \, ds.
\]
Recalling that the mapping \(\bar{\theta} \mapsto \theta\) is Lipschitz continuous from \(L^\infty(0, T; L^p(\Omega))\) to \(L^\infty(0, T; L^p(\Omega))\) and maps any bounded subset of \(L^\infty(0, T; L^p(\Omega))\) into a bounded subset of \(L^\infty(0, T; L^p(\Omega))\) with \(p = 2\) and \(\bar{q} > 4\), \(p \geq 4\), \(q = \beta_1 \bar{q} > 8\), the last estimate allows us to conclude.

Let us observe furthermore that the image of a bounded set of \(L^\infty(0, T; L^p(\Omega))\) by the mapping \(\bar{\theta} \mapsto \zeta(\bar{\theta}) = \theta \mapsto (u, z)\) is a bounded subset of \(H^1(0, T; H^1_0(\Omega) \times L^2(\Omega)) \cap L^\infty(0, T; H^1_0(\Omega) \times H^1(\Omega))\).

Let us introduce now some new notations. For any \(r > 1\), let
\[
V(\Omega; \mathbb{R}^3) \overset{\text{def}}{=} \{u \in L^2(\Omega; \mathbb{R}^3) : \nabla u \in L^r(\Omega; \mathbb{R}^{3 \times 3})\},
\]
and for any \(r \geq 2\), let
\[
V_0(\Omega; \mathbb{R}^3) \overset{\text{def}}{=} \{u \in V(\Omega; \mathbb{R}^3) : u_{\mid_{\partial\Omega}} = 0\}.
\]
We endowed \( V^r(\Omega; \mathbb{R}^3) \) with the following norm
\[
\forall u \in V^r(\Omega; \mathbb{R}^3) : \| u \|_{V^r(\Omega)} \overset{\text{def}}{=} \| u \|_{L^2(\Omega)} + \| \nabla u \|_{L^r(\Omega)}.
\]

The aim of the next two lemmas is to prove further regularity results for the solutions \((u, z)\) of the system composed by the momentum equilibrium equation and the flow rule. More precisely, assuming that \( \theta \) remains in a bounded subset of \( L^q(0,T;L^p(\Omega)) \), we will prove that \( e(\dot{u}), \dot{z} \) and \( z \) remain in a bounded subset of \( L^q(0,T;\mathbb{R}_+^N) \cap L^q/2(0,T;L^p(\Omega)) \) and \( L^q(0,T;H^2(\Omega)) \), respectively.

**Lemma 4.4** Assume that (3.10), (3.11), (3.13b), (3.14) and (3.10) hold. Assume moreover that \( z^0 \in X_{q,p}(\Omega) \) and \( u^0 \in V^r_0(\Omega; \mathbb{R}^3) \). Then \( e(u) \) belongs to \( W^{1,q}(0,T;L^p(\Omega)) \) and \( \theta \mapsto e(u) \) maps any bounded subset of \( L^q(0,T;L^p(\Omega)) \) into a bounded subset of \( W^{1,q}(0,T;L^p(\Omega)) \).

**Proof.** The idea of the proof is to interpret (4.1a) as an ODE for \( \phi \) and \( \Lambda \).

The aim of the next two lemmas is to prove further regularity results for the solutions \((u, z)\), such that
\[
\forall (\Omega), \text{with notations:}
\]
\[ (\Omega) \]

The proof follows that \((4.1a)\) can be rewritten as
\[
\| u \|_{V^r(\Omega;\mathbb{R}^3)} \leq C(\| \phi_1 \|_{L^2(\Omega)} + \| \phi_2 \|_{L^r(\Omega)}) = C\| \phi \|_{\mathcal{F}_p},
\]
and \( \Lambda_p \) is linear continuous from \( \mathcal{F}_p \) to \( V^r_0(\Omega; \mathbb{R}^3) \) (for more details, the reader is referred to [Val88]). It follows that \((1.1a)\) can be rewritten as
\[
\dot{u} = \mathcal{G}_p(\varphi_{\theta} u),
\]
where \( \varphi_{\theta} \overset{\text{def}}{=} (\ell, E z - \alpha \theta I) \) and \( \mathcal{G}_p(\varphi_{\theta} u) \overset{\text{def}}{=} \Lambda_p(\varphi_{\theta} - \mathcal{A}_p u) = \Lambda_p(\varphi_{\theta} - \mathcal{A}_p(\mathcal{A}_p u)). \)

With the assumption (3.10) and the previous results, we already know that \( \varphi_{\theta} \in L^q(0,T;\mathcal{F}_p) \) and we can apply classical results for ODE in Banach spaces to conclude that \( u \in W^{1,q}(0,T;V^r_0(\Omega; \mathbb{R}^3)) \), the reader is referred to [Car90] for more details.

We can also obtain estimates for \( u \) and \( \Lambda_p \in V^r_0(\Omega; \mathbb{R}^3) \). To this aim, we introduce the following notations: \( \lambda_{\Lambda} \overset{\text{def}}{=} \| \Lambda_p \|_{L(\mathcal{F}_p, V^r_0(\Omega))} \) and \( \lambda_{\mathcal{A}} \overset{\text{def}}{=} \| \mathcal{A}_p \|_{L(V^r_0(\Omega), \mathcal{F}_p)} \).

Then, we observe that \((1.41)\) gives
\[
\| \dot{u}(\cdot,t) \|_{V^r(\Omega)} \leq C_{\lambda_{\Lambda}} (\| \varphi_{\theta}(\cdot,t) - \mathcal{A}_p u_0 \|_{\mathcal{F}_p} + C_{\lambda_{\mathcal{A}}} \| u(\cdot,t) - u_0 \|_{V^r(\Omega)}) \text{ a.e. } t \in [0,T].
\]

Let us turn now to the term \( \| u(\cdot,t) - u_0 \|_{V^r(\Omega)} \). It is clear that
\[
\| u(\cdot,t) - u_0 \|_{V^r(\Omega)} \leq \int_0^t \| \mathcal{G}_p(\varphi_{\theta}(\cdot,s), u(\cdot,s)) \|_{V^r(\Omega)} ds
\]
\[
\leq C_{\lambda_{\Lambda}} \int_0^t \| \varphi_{\theta}(\cdot,s) - \mathcal{A}_p u_0 \|_{\mathcal{F}_p} ds + C_{\lambda_{\Lambda}} C_{\lambda_{\mathcal{A}}} \int_0^t \| u(\cdot,s) - u_0 \|_{V^r(\Omega)} ds.
\]

Therefore, we may infer from Grönwall’s lemma and the continuous embedding \( H^1(\Omega) \hookrightarrow L^p(\Omega) \) that there exists a generic constant \( C_1 > 0 \) such that
\[
\| u(\cdot,t) - u_0 \|_{V^r(\Omega)} \leq e^{C_{\lambda_{\Lambda}} C_{\lambda_{\mathcal{A}}} t} C_{\lambda_{\Lambda}} \int_0^t \| \varphi_{\theta}(s) - \mathcal{A}_p u_0 \|_{\mathcal{F}_p} ds
\]
\[
\leq e^{C_{\lambda_{\Lambda}} C_{\lambda_{\mathcal{A}}} C_1} C_{\lambda_{\Lambda}} C_1 \int_0^t [\| \ell(s) \|_{L^2(\Omega)} + \| \mathcal{E} \|_{L^\infty(\Omega)} + \| z(\cdot,s) \|_{H^1(\Omega)} + \| \theta(\cdot,s) \|_{L^p(\Omega)} + C_{\lambda_{\mathcal{A}}} \| u_0 \|_{V^r(\Omega)}] ds.
\]
We insert (4.43) in (4.42), we find that there exists $C_2 > 0$ such that
\[
\|\bar{\alpha}\|_{L^\infty(0,T;V^\prime(\Omega))} \leq C_2 (\|\|L^\infty(\Omega)\| + \|\|H^1(\Omega)\| + \|\|\|u^0\|_{V^\prime(\Omega)}\|).
\]
Recalling Lemma 4.3 we infer that the image of any bounded set of $L^2(0,T;L^2(\Omega))$ by the mappings \(\bar{\alpha} \mapsto \zeta(\bar{\alpha}) = \theta \mapsto e(u)\) and \(\bar{\alpha} \mapsto \zeta(\bar{\alpha}) = \theta \mapsto e(\hat{u})\) are still bounded sets in $L^\infty(0,T;L^p(\Omega))$ and $L^2(0,T;L^p(\Omega))$, respectively.

Let us conclude this section with some regularity results for $z$ and $\dot{z}$.

**Lemma 4.5** Assume that (3.10), (3.1b), (3.13a), (3.13c) and (3.1d) hold. Assume moreover that $z^0 \in X_{q,p}(\Omega)$ and $u^0 \in V^p_0(\Omega;\mathbb{R}^3)$. Then $\dot{z}$ and $\Delta z$ belong to $L^{q/2}(0,T;L^q(\Omega)) \cap L^q(0,T;L^{q/2}(\Omega))$ and $z \in C^0([0,T],X_{q,p}(\Omega)) \cap L^q(0,T;H^2(\Omega))$. Moreover $\theta \mapsto (\dot{z},\Delta z,z)$ maps any bounded subset of $L^q(0,T;L^2(\Omega))$ into a bounded subset of $L^q(0,T;L^p(\Omega)) \cap L^q(0,T;L^{p/2}(\Omega))^2 \times (C^0([0,T],X_{q,p}(\Omega)) \cap L^q(0,T;H^2(\Omega)))$.

**Proof.** Let us rewrite again (4.1b) as
\[
\dot{z} - \nu M^{-1} \Delta z = M^{-1} f^z.
\]
with $f^z \triangleq \text{Proj}_{3 \times 3}(\mathbb{E}(e(u) - z)) - D_z H_i(z) \theta D_z H_2(z) - \psi$ and $\psi \triangleq \text{Proj}_{3 \times 3}(w) - M \dot{z} \in \partial \Psi_{\text{dev}}(\dot{z})$ where $w$ has been defined in the proof of Theorem 4.1. With the assumption (3.10a), we have
\[
\forall \psi \in \partial \Psi_{\text{dev}}(\dot{z}) : \|\psi\| \leq C\Psi \text{ a.e. } (x,t) \in \Omega \times (0,T).
\]
Then Lemma 4.3 enables us to infer that $\text{Proj}_{3 \times 3}(\mathbb{E}(e(u) - z)) - \Psi$ remains bounded in $L^\infty(0,T;L^p(\Omega))$. Furthermore, we know with Lemma 4.3 that $z$ is bounded in $L^\infty(0,T;\mathbb{H}^1(\Omega))$, so, using 3.12, we infer that $D_z H_i(z)$ is bounded in $L^\infty(0,T;L^p(\Omega))$ for $i = 1,2$. Then it follows that $M^{-1} f^z$ is bounded in $L^q(0,T;L^p(\Omega))$ if $\theta$ belongs to a bounded subset of $L^q(0,T;L^p(\Omega))$. We may deduce from the maximal regularity result for parabolic systems that $z$ is bounded in $L^q(0,T;H^2(\Omega))$ and $\dot{z}$ is bounded in $L^q(0,T;L^{p/2}(\Omega))$ for $i = 1,2$.

Since $H^2(\Omega) \hookrightarrow L^\infty(\Omega)$ with continuous embedding, $z$ is bounded in $L^q(0,T;L^\infty(\Omega))$ and thus $\theta D_z H_2(z)$ is bounded in $L^q(0,T;L^p(\Omega))$. We may deduce that $f^z$ is bounded in $L^q(0,T;L^p(\Omega))$ and, the maximal regularity result for parabolic systems allows us to infer that $\dot{z}$ and $\Delta z$ belong to a bounded subset of $L^q(0,T;L^{p/2}(\Omega)) \cap L^q(0,T;L^p(\Omega))$ and $z$ belongs to a bounded subset of $C^0([0,T],X_{q,p}(\Omega))$ whenever $\theta$ belongs to a bounded subset of $L^q(0,T;L^p(\Omega))$. \qed

## 5 Existence and regularity results for the enthalpy equation

In this section existence and uniqueness results for the enthalpy equation are recalled and some regularity results are obtained. More precisely, let us consider the enthalpy equation $(P_o)$:
\[
\dot{\vartheta} - \text{div}(\kappa^c \nabla \vartheta) = f
\]
with initial conditions
\[
\vartheta(0) = \vartheta^0,
\]
and boundary conditions
\[
\kappa^c \nabla \vartheta \cdot n_{\partial \Omega} = 0.
\]
We assume that the initial enthalpy $\vartheta^0$ belongs to $L^2(\Omega)$ and $f$ belongs to $L^2(0,T;L^2(\Omega))$. Furthermore, we assume that $\kappa^c \in L^\infty(\mathbb{Q}_T;\mathbb{R}^{3 \times 3}_{\text{sym}})$ and satisfies
\[
\exists c_{\kappa^c} > 0 \forall v \in \mathbb{R}^3 : \kappa^c(x,t) v \cdot v \geq c_{\kappa^c} |v|^2 \text{ a.e. } (x,t) \in \mathbb{Q}_T,
\]
\[
\exists c_{\kappa^c} > 0 : |\kappa^c(x,t)| \leq C_{\kappa^c} \text{ a.e. } (x,t) \in \mathbb{Q}_T.
\]
The weak formulation of the problem is given by
\[
\begin{cases}
\text{Find } \vartheta : [0,T] \to H^1(\Omega) \text{ such that } \vartheta(0) = \vartheta^0 \text{ and for all } \xi \in H^1(\Omega),
\int_{\Omega} \dot{\vartheta} \xi dx + \int_{\Omega} \kappa^c \nabla \vartheta \cdot \nabla \xi dx = \int_{\Omega} f \xi dx \text{ in the sense of distributions.}
\end{cases}
\]
Theorem 5.1 (Existence and uniqueness for \((P_\theta)\)) Under the previous assumptions, the problem \(\text{(5.1) - (5.3)}\) possesses a unique solution \(\vartheta \in C^0([0,T];L^2(\Omega) \cap L^2(0,T;H^1(\Omega)))\) with \(\dot{\vartheta} \in L^2(0,T; (H^1(\Omega))').\) Moreover we have
\[
\forall \tau \in [0,T] : \|\vartheta(\tau)\|_{L^2(\Omega)}^2 + 2\kappa \int_0^\tau \|\nabla \dot{\vartheta}(t)\|_{L^2(\Omega)}^2 \, dt \leq e^\kappa \left(\|\vartheta^0\|_{L^2(\Omega)}^2 + \|f\|_{L^2(0,T;L^2(\Omega))}\right).
\]

Proof. The proof of existence and uniqueness of a solution is quite classical and can be found in [Bre83], [Ra83]. The estimate is straightforward and its verification is left to the reader. \(\square\)

Let us introduce the following functional space
\[
\mathcal{W} \overset{\text{def}}{=} \{\vartheta \in L^2(0,T; H^1(\Omega)) \cap L^\infty(0,T; L^2(\Omega)) : \dot{\vartheta} \in L^2(0,T; (H^1(\Omega))')\},
\]
endowed with the norm
\[
\forall \vartheta \in \mathcal{W} : \|\vartheta\|_\mathcal{W} \overset{\text{def}}{=} \|\vartheta\|_{L^2(0,T; H^1(\Omega))} + \|\vartheta\|_{L^\infty(0,T; L^2(\Omega))} + \|\dot{\vartheta}\|_{L^2(0,T; (H^1(\Omega))')}.
\]
Due to [Sim87], we know that \(\mathcal{W}\) is compactly embedded in \(L^q(0,T; L^\varphi(\Omega))\). Note that the previous estimate implies that there exists a generic constant \(C > 0\) such that the solution of problem \((P_\vartheta)\) satisfies
\[
\|\vartheta\|_\mathcal{W} \leq C\left(\|\vartheta^0\|_{L^2(\Omega)} + \|f\|_{L^2(0,T; L^2(\Omega))}\right).
\]

6 Local existence result

This section is dedicated to the proof of a local existence result for \(\text{(3.4) - (3.8)}\) by using a fixed-point argument. To this aim, for any given \(\vartheta \in L^\varphi(0,T; L^\varphi(\Omega))\), we consider \(\hat{\kappa}_c \overset{\text{def}}{=} \kappa^c(e(u), z, \theta)\) and \(f = f^\vartheta = Lc(\dot{\vartheta}) + \theta(\alpha \text{tr}(e(\dot{\vartheta}))) + D_z H_2(x; \dot{z}) + \Psi(\dot{z}) + M \dot{z}; \dot{z}\) in \((P_\vartheta)\), where \((u, z)\) are the solutions of \((P_{uz})\) with \(\theta = \zeta(\vartheta)\). With the results obtained in Section 4 we already know that \(f^\vartheta\) belongs to \(L^q(0,T; L^{\varphi/2}(\Omega))\). Since \(p \geq 4\) and \(q > 8\), we infer that \(f^\vartheta\) belongs to \(L^2(0,T; L^2(\Omega))\) and we can define \(\vartheta \in C^0([0,T]; L^2(\Omega)) \cap \mathcal{W}\) as the unique solution of \((P_\vartheta)\). Thus we can introduce the fixed point mapping \(\phi : \vartheta \mapsto \vartheta\) from \(L^\varphi(0,T; L^\varphi(\Omega))\) to \(L^\varphi(0,T; L^\varphi(\Omega))\).

Proposition 6.1 The mapping \(\phi\) is continuous from \(L^\varphi(0,T; L^\varphi(\Omega))\) to \(L^\varphi(0,T; L^\varphi(\Omega))\).

Proof. Let \((\vartheta_n)_{n\in\mathbb{N}}\) be a converging sequence of \(L^\varphi(0,T; L^\varphi(\Omega))\) and let \(\vartheta_s\) be its limit. We denote by \(\vartheta_n \overset{\text{def}}{=} \phi(\vartheta_n)\) for all \(n \geq 0\) and \(\vartheta_s \overset{\text{def}}{=} \phi(\vartheta_s)\). Since \((\vartheta_n)_{n\in\mathbb{N}}\) is a bounded family of \(L^\varphi(0,T; L^\varphi(\Omega))\), we infer from the previous results that \((\vartheta_n)_{n\in\mathbb{N}}\) is bounded in \(C^0([0,T]; L^2(\Omega)) \cap \mathcal{W}\). We may deduce that \((\vartheta_n)_{n\in\mathbb{N}}\) is relatively compact in \(L^2(0,T; L^2(\Omega))\) (see [Sim87]) and there exists a subsequence, still denoted \((\vartheta_n)_{n\in\mathbb{N}}\), such that
\[
\vartheta_n \rightharpoonup \vartheta\quad \text{in} \quad L^2(0,T; H^1(\Omega)) \quad \text{weak},
\]
\[
\vartheta_n \rightarrow \vartheta\quad \text{in} \quad L^\varphi(0,T; L^\varphi(\Omega)).
\]
Let us define \(\mathcal{V}_T \overset{\text{def}}{=} \{w \in C^\infty([0,T]) : w(T) = 0\}\). Hence we observe that for all \(n \geq 0\), we have
\[
\forall \xi \in H^1(\Omega) \forall w \in \mathcal{V}_T : - \int_{Q_T} \vartheta_n(x,t)\xi(x)w(t) \, dx \, dt + \int_{Q_T} \hat{\kappa}_c \nabla \vartheta_n(x,t) \nabla \xi(x)w(t) \, dx \, dt
\]
\[
= \int_{Q_T} f^\vartheta_n(x,t)\xi(x)w(t) \, dx \, dt + \int_{Q_T} \vartheta^0(x)\xi(x)w(0) \, dx,
\]
with \(\hat{\kappa}_c \overset{\text{def}}{=} \kappa^c(e(u_n), z_n, \theta_n)\) and \((u_n, z_n)\) solutions of \((P_{uz})\) with \(\theta_n \overset{\text{def}}{=} \zeta(\vartheta_n)\). Since \((\vartheta_n)_{n\in\mathbb{N}}\) converges to \(\vartheta_s\) in \(L^\varphi(0,T; L^\varphi(\Omega))\), we infer from Lemma 4.3 that \((u_n, z_n)_{n\in\mathbb{N}}\) converges to the solution \((u_s, z_s)\) of \((P_{uz})\) with \(\theta_s = \zeta(\vartheta_s)\) in \(H^1(0,T; H_0^1(\Omega) \cap L^\infty(0,T; H^1(\Omega)))\) and \(\vartheta^0(x)\xi(x)w(0)\) \(\mathcal{V}_T\). Let us recall that the mapping \(\phi_1 : \vartheta \mapsto \theta = \zeta(\vartheta)\) is Lipschitz continuous from \(L^\varphi(0,T; L^\varphi(\Omega))\) to \(L^\varphi(0,T; L^\varphi(\Omega))\), which implies, possibly after extracting another subsequence, that
\[
\theta_n, u_n, z_n \rightarrow \theta_s, u_s, z_s \quad \text{a.e.} \quad (x,t) \in Q_T.
\]
Note that the continuity of the mapping \( \kappa^c \) gives
\[
\tilde{\kappa}^c_n = \kappa^c(e(u_n), z_n, \theta_n) \to \tilde{\kappa}^c_s = \kappa^c(e(u_s), z_s, \theta_s) \text{ a.e. in } Q_T,
\]
and due to the boundedness assumption on \( \kappa^c \), we obtain with the Lebesgue’s theorem that
\[
\tilde{\kappa}^c_n \nabla \xi w \to \tilde{\kappa}^c_s \nabla \xi w \text{ in } L^2(0, T; L^2(\Omega)).
\]
Therefore it is possible to pass to the limit in all the terms of the left hand side of (6.11) to get
\[
\forall \xi \in H^1(\Omega) \forall w \in V_T : - \int_{Q_T} \vartheta(x, t) \xi(x) \dot{w}(t) \, dx \, dt + \int_{Q_T} \tilde{\kappa}^c_n \nabla \vartheta(x, t) \nabla \xi(x) \, w(t) \, dx \, dt = \lim_{n \to +\infty} \int_{Q_T} f^{\tilde{\kappa}^c_n}(x, t) \xi(x) \, w(t) \, dx \, dt + \int_{\Omega} \vartheta^0(x) \xi(x) \, w(0) \, dx.
\]
Recalling that \( p \in [4, 6] \) and \( q > 8 \), we infer that the mapping \( \vartheta \mapsto f^{\tilde{\vartheta}} \) is continuous from \( L^{\tilde{\vartheta}}(0, T; L^p(\Omega)) \) to \( L^{r_s}(0, T; L^2(\Omega)) \) with \( \frac{1}{r_1} = \frac{1}{q} + \frac{1}{2} \) and \( \frac{1}{r_2} = \frac{1}{p} + \frac{1}{2} \). Indeed, for any \( \vartheta_i \) in \( L^{\tilde{\vartheta}}(0, T; L^p(\Omega)) \), let \( (u_i, z_i) \) be the solution of \( (P_{u_s}) \) with \( \vartheta_i = \zeta(\vartheta_i) \), \( i = 1, 2 \), we find
\[
\begin{align*}
f^{\tilde{\vartheta}_1} - f^{\tilde{\vartheta}_2} &= \frac{1}{2} \int_{\Omega} (\dot{u}_1 + \dot{u}_2) : e(\dot{u}_1 - \dot{u}_2) \, dx + \left( \frac{1}{\theta_1} - \frac{1}{\theta_2} \right) \alpha \text{tr}(e(\dot{u}_1) + D_2 H_2(z_1) : \dot{z}_1) + \frac{1}{2} \alpha \text{tr}(e(\dot{u}_1 - \dot{u}_2) + D_2 H_2(z_1) : \dot{z}_1 - D_2 H_2(z_2) : \dot{z}_2) + \Psi(\dot{z}_1) - \Psi(\dot{z}_2) + M(\dot{z}_1 + \dot{z}_2) : (\dot{z}_1 - \dot{z}_2).\end{align*}
\]
On the other hand, (3.10a) and (5.10a) imply
\[
|\Psi(\dot{z}_1) - \Psi(\dot{z}_2)| \leq C^\Psi |\dot{z}_1 - \dot{z}_2|,
\]
and (5.10b) and (3.11b) give
\[
|D_2 H_2(z_1) : \dot{z}_1 - D_2 H_2(z_2) : \dot{z}_2| \leq |D_2 H_2(z_1)||\dot{z}_1 - \dot{z}_2| + |D_2 H_2(z_2) - D_2 H_2(z_2)||\dot{z}_2| \leq C_{H_2} |1 + |z_1||\dot{z}_1 - \dot{z}_2| + C_{H_2} |z_1 - z_2||\dot{z}_2|.
\]
The boundedness and the continuity properties proved in Lemma 4.4 and Lemma 4.3, respectively, allow us to deduce the desired result. Therefore, we may infer that
\[
\forall \xi \in H^1(\Omega) \forall w \in V_T : \lim_{n \to +\infty} \int_{Q_T} f^{\tilde{\kappa}^c_n}(x, t) \xi(x) \, w(t) \, dx \, dt = \int_{Q_T} f^{\tilde{\kappa}^c_s}(x, t) \xi(x) \, w(t) \, dx \, dt.
\]
We conclude that \( \vartheta \) is solution of problem \( (P_{u_s}) \) with the data \( \tilde{\kappa}^c_s \) and \( f^{\tilde{\vartheta}_s} \). Moreover by uniqueness of the solution, it follows that \( \vartheta = \vartheta_\ast \) and the whole sequence \( (\vartheta_n)_{n \in N} \) converges to \( \vartheta_\ast = \zeta(\vartheta_\ast) \). \( \square \)

We establish now that the mapping \( \varphi \) fulfills the other assumptions of the Schauder’s fixed point theorem. To this aim, we introduce some notations: let \( R^0, R^\varphi > 0 \) be any given positive real numbers such that \( \max(\|u^0\|_{V_T(\Omega)}, \|z^0\|_{X_{\alpha_1, \alpha_2}(\Omega)}) \leq R^0 \) and \( \|\vartheta\|_{L^6(0, T; L^p(\Omega))] \leq R^\varphi \). Clearly, we have
\[
\|\zeta(\vartheta)\|_{L^6(0, T; L^p(\Omega))] \leq \|\vartheta\|_{L^6(0, T; L^p(\Omega))] \leq (\frac{\|\vartheta\|_{L^6(0, T; L^p(\Omega))]}}{R^\varphi}) \|\vartheta\|_{L^6(0, T; L^p(\Omega))] \leq R^\varphi \text{ def } (\frac{\|\vartheta\|_{L^6(0, T; L^p(\Omega))]}}{R^\varphi}) \frac{\|\vartheta\|_{L^6(0, T; L^p(\Omega))]}}{R^\varphi}.
\]
Therefore once again the results of Section 9 are used, which imply that there exists a constant \( R^\varphi \text{ def } R^l(R^0, R^\varphi, \|\vartheta\|_{L^6(0, T; L^p(\Omega))]}) > 0 \), depending only on \( R^0 \), \( R^\varphi \) and \( \|\vartheta\|_{L^6(0, T; L^p(\Omega))] \), such that
\[
\|f^{\tilde{\vartheta}}\|_{L^{6/4}(0, T; L^{p/2}(\Omega))] \leq R^l(R^0, R^\varphi, \|\vartheta\|_{L^6(0, T; L^p(\Omega))]}. \]

The results of Section 9 imply that there exists a generic constant \( C > 0 \) such that
\[
\|\vartheta\|_{L^\infty(0, T; L^2(\Omega))] \leq C\|f^{\tilde{\vartheta}}\|_{L^2(0, T; L^2(\Omega))} + \|\vartheta^0\|_{L^2(\Omega)] \leq C\|\vartheta\|_{L^\infty(0, T; L^2(\Omega))] \leq C\|f^{\tilde{\vartheta}}\|_{L^{6/4}(0, T; L^{p/2}(\Omega))] + \|\vartheta^0\|_{L^2(\Omega)] \leq C\|\vartheta\|_{L^\infty(0, T; L^2(\Omega))] \leq C\|f^{\tilde{\vartheta}}\|_{L^{6/4}(0, T; L^{p/2}(\Omega))] + \|\vartheta^0\|_{L^2(\Omega)].
\]
Now let $0 < \tau \leq T$ and let us introduce the following functional space

$$\mathcal{W}_\tau \overset{df}{=} \{ \vartheta \in L^2(0, \tau; \mathcal{H}^1(\Omega)) \cap L^\infty(0, \tau; L^2(\Omega)) : \dot{\vartheta} \in L^2(0, \tau; (\mathcal{H}^1(\Omega))') \}.$$ 

For any $\bar{\vartheta} \in L^q(0, \tau; L^p(\Omega))$, we define its extension $\tilde{\vartheta}_{\text{ext}}$ by $\tilde{\vartheta}_{\text{ext}} = \bar{\vartheta}$ on $[0, \tau]$ and $\tilde{\vartheta}_{\text{ext}} = 0$ on $(\tau, T]$. It is clear that $\tilde{\vartheta}_{\text{ext}} \in L^q(0, T; L^p(\Omega))$ and the mapping $\bar{\vartheta} \mapsto \tilde{\vartheta}_{\text{ext}}$ is a contraction from $L^q(0, \tau; L^p(\Omega))$ into $L^q(0, T; L^p(\Omega))$. For any $\tilde{\vartheta} \in L^q(0, \tau; L^p(\Omega))$, we define $\phi_{\tau}(\tilde{\vartheta})$ as the restriction on $[0, \tau]$ of $\tilde{\vartheta}_{\text{ext}}$.

We infer immediately from Proposition 6.1 that $\phi_{\tau}$ is continuous from $L^q(0, \tau; L^p(\Omega))$ to $L^q(0, \tau; L^p(\Omega))$. Furthermore, for any $\tilde{\vartheta} \in L^q(0, \tau; L^p(\Omega))$, we have

$$\|\phi_{\tau}(\tilde{\vartheta})\|_{L^q(0, \tau; L^p(\Omega))} = \|\phi_{\tau}(\tilde{\vartheta})\|_{L^q(0, \tau; L^p(\Omega))} = \left( \int_0^\tau \|\phi(\tilde{\vartheta}_{ext}(\cdot, t))\|^q_{L^2(\Omega)} \, dt \right)^{\frac{1}{q}} \leq T^{\frac{1}{q}} \|\phi(\tilde{\vartheta}_{ext})\|_{L^\infty(0, T; L^2(\Omega))},$$

and the previous estimates allow us to show that, for any $R^d > 0$, there exists $\tau \in (0, T]$ such that $\phi_{\tau}$ maps the closed ball $B_{L^q(0, \tau; L^p(\Omega))}(0, R^d)$ into itself. Note that the image of $B_{L^q(0, \tau; L^p(\Omega))}(0, R^d)$ by $\phi$ is a bounded subset of $\mathcal{W}$ and thus it is relatively compact in $L^q(0, T; L^p(\Omega))$. It follows that the image of $B_{L^q(0, \tau; L^p(\Omega))}(0, R^d)$ by $\phi_{\tau}$ is also relatively compact in $L^q(0, \tau; L^p(\Omega))$. Consequently, we may conclude that the problem (3.6)–(3.8) possesses a local solution $(u, z, \vartheta)$ defined on $[0, \tau]$ such that $u \in W^1(0, \tau; V_0^p(\Omega)), z \in L^\infty(0, \tau; \mathcal{H}^1(\Omega)) \cap H^1(0, \tau; L^2(\Omega)) \cap C^0(0, \tau; X_{q,p}(\Omega)) \cap L^q(0, \tau; H^2(\Omega)), \dot{z}, \Delta z \in L^\infty(0, \tau; L^2(\Omega)) \cap L^q(0, \tau; L^p(\Omega))$ and $\vartheta \in \mathcal{W}_\tau$.

We have to go back to the problem (3.1)–(3.3). First we observe that $g$ and $\zeta$ define a $C^1$-diffeomorphism from $(0, \infty)$ to $(0, \infty)$ and any solution of (3.6)–(3.8) provides a solution of (3.1)–(3.3) as soon as the enthalpy $\vartheta$ remains strictly positive. So we assume now that the initial enthalpy is strictly positive almost everywhere on $\Omega$, i.e., there exists $\bar{\vartheta} > 0$ such that

$$g(\theta^0(x)) = \vartheta^0(x) \geq \bar{\vartheta} > 0 \text{ a.e. } x \in \Omega.$$

Therefore it is possible to use the Stampacchia’s truncation method and to prove a local existence result for the problem (3.1)–(3.3).

**Theorem 6.2 (Local existence result)** Assume that (3.10), (3.11), (3.13), (3.14), (3.16) and (3.17) hold. Then, for any initial data $u^0 \in V_0^p(\Omega), z^0 \in X_{q,p}(\Omega)$ and $\vartheta^0 \in L^2(\Omega)$ satisfying (3.2), there exists $\tau \in (0, T]$ such that the problem (3.1)–(3.3) admits a solution on $[0, \tau]$.

**Proof.** Let $(u, z, \vartheta)$ be a solution of (3.6)–(3.8) on $[0, \tau]$. We prove now that

$$\vartheta(x, t) > 0 \text{ a.e. } (x, t) \in Q_\tau.$$

To do so, we introduce some notations: let $C \overset{df}{=} (\lambda_0)^2 + (\mathcal{H}_2)^2 \nu$ and let $\varphi : [0, \tau] \to \mathbb{R}$ such that

$$\forall t \in [0, \tau] : \varphi(t) \overset{df}{=} \frac{\dot{\vartheta}}{\vartheta} - \Delta \int_0^t \left[ C + \frac{\mathcal{H}_2\nu}{\vartheta}(z(s), s) \right] \|z\|^2_{L^\infty(\Omega)} \, ds.$$

We use the classical Stampacchia’s truncation method. Let us define $G \in C^1(\mathbb{R})$ satisfying

(i) $\forall \sigma \in \mathbb{R} \exists C^G > 0 : |G'(\sigma)| \leq C^G$,

(ii) $G$ is strictly increasing on $(0, \infty)$,

(iii) $\forall \sigma \leq 0 : G(\sigma) = 0$.

Let us define also $H(\sigma) \overset{df}{=} \int_0^\sigma G(s) \, ds$ for all $\sigma \in \mathbb{R}$, $\vartheta_1 \overset{df}{=} -\vartheta + \varphi$ and $h(t) \overset{df}{=} \int_0^t H(\vartheta_1) \, dx$. Clearly, we have $H \in C^2(R; \mathbb{R})$ and $H(\sigma) > 0$ for all $\sigma > 0$. Furthermore, $\vartheta_1(0) = -\vartheta^0 + \bar{\vartheta} \leq 0$ almost everywhere on $\Omega$ implies that $h(0) = 0$. Since $\vartheta \in \mathcal{W}_\tau$ and $\varphi \in H^1(0, \tau; \mathbb{R})$, we infer that $h$ is absolutely continuous and

$$\dot{h}(t) = \int_\Omega G(\vartheta_1) \vartheta_1 \, dx = -\int_\Omega G(\vartheta_1) \div (\kappa \epsilon \nabla \vartheta_1) + \langle \mathbf{L}(\dot{u}):e(\dot{u}) + \theta(\alpha |\epsilon(e(\dot{u}))| + D_z H_2(z):\dot{z} + \Psi(\dot{z}) + M \dot{z}:\dot{z} - \dot{\varphi} \rangle \, dx = -\int_\Omega G'(\vartheta_1) \kappa \overset{\circ}{\nabla} \vartheta_1 : \overset{\circ}{\nabla} \vartheta_1 \, dx - \int_\Omega \varphi(\dot{u}) \div (\mathbf{L}(\dot{u}):e(\dot{u}) + \theta(\alpha |\epsilon(e(\dot{u}))| + D_z H_2(z):\dot{z} + \Psi(\dot{z}) + M \dot{z}:\dot{z} - \dot{\varphi} \rangle \, dx,$$
for almost every $t \in [0, \tau]$. It follows from (3.12), (3.14) and Cauchy-Schwarz’s inequality that

$$\text{Le}(\dot{u}):e(\dot{u}) + \alpha \theta \text{tr}(e(\dot{u})) \geq c^\theta |e(\dot{u})|^2 - 3\alpha |\theta||e(\dot{u})| \geq \frac{c^\theta}{2}|e(\dot{u})|^2 - \frac{3(\alpha)^2|\theta|^2}{2c^\theta},$$

and

$$Mz:z + \theta D_z H_2(z):z \geq c^M |z|^2 - |\theta||D_z H_2(z):z| \geq c^M |z|^2 - C^H_2|\theta|(1+|z|)|z| \geq \frac{c^M}{2}|z|^2 - \frac{C^H_2|\theta|^2}{c^\theta}(1+|z|^2).$$

Since $G'(\vartheta_1) \geq 0$ and $G(\vartheta_1) \geq 0$ almost everywhere and (3.10) and (3.17) hold, we get

$$\dot{h}(t) \leq \int \Omega G(\vartheta_1)\left(|\theta|^2\left(C + \frac{C^H_2|\theta|^2}{c^\theta}|z|^2\right) + \varphi\right)dx \text{ a.e. } t \in [0, \tau].$$

But $\theta = \zeta(\vartheta)$, and reminding that $\beta_1 \geq 2$, we have with (3.18) that

$$|\theta| = |\zeta(\vartheta)| \leq \sqrt{\frac{2\varphi + \varphi}{\varphi} + 1 - 1} \leq \sqrt{\frac{2\varphi}{\varphi}} \text{ a.e. } (x, t) \in Q_{\tau}.$$

On the other hand, $G(\vartheta_1)$ vanishes whenever $\vartheta \geq \varphi$, it follows that

$$\dot{h}(t) \leq \int \Omega G(\vartheta_1)\left(\frac{\varphi}{e(\vartheta)}\varphi\left(C + \frac{C^H_2|\theta|^2}{c^\theta}|z|^2\right) + \varphi\right)dx \leq 0 \text{ a.e. } t \in [0, \tau].$$

We may deduce that $h(t) \leq h(\tilde{t}) = 0$ for all $t \in [0, \tau]$. Then we infer that

$$H(\vartheta_1) = 0 \text{ a.e. } (x, t) \in \Omega \times (0, \tau),$$

which implies that

$$\vartheta_1 = -\vartheta + \varphi \leq 0 \text{ a.e. } (x, t) \in \Omega \times (0, \tau).$$

This concludes the proof. \(\square\)

7 Global existence result

We begin this section with some a priori estimates for the solutions of the problem (3.0)–(3.8). As usual, the result relies on an energy balance combined with Grönnwalt’s lemma. Then the global existence result is proved by using a contradiction argument together with the results obtained in the previous sections.

**Proposition 7.1 (Global energy estimate)** Assume that (3.10), (3.11), (3.13), (3.14), (3.16) and (3.17) hold. Assume moreover that $u^0 \in V_Q^0(\Omega)$, $\vartheta^0 \in X_{u,p}(\Omega)$, $\theta^0 \in L^2(\Omega)$ such that (6.2) holds. Then, there exists a constant $\tilde{C} > 0$, depending only on $\|u^0\|^2_{H^1(\Omega)}$, $\|z^0\|^2_{H^1(\Omega)}$, $\|\vartheta^0\|_{L^1(\Omega)}$ and the data such that for any solution $(u, z, \vartheta)$ of problem (3.0)–(3.8) defined on $[0, \tau]$, $\tau \in (0, T]$, we have

$$\forall \tau \in [0, \tau] : \|u(\cdot, \tau)\|^2_{H^1(\Omega)} + \|z(\cdot, \tau)\|^2_{H^1(\Omega)} + \|\vartheta(\cdot, \tau)\|_{L^1(\Omega)} \leq \tilde{C}.$$

**Proof.** On the one hand, we multiply (3.6b) by $\dot{u}$ and we integrate this expression over $Q_{\tau}$, with $\tau \in [0, \tau]$, to get

$$\int_{Q_{\tau}} \left(\widehat{\Theta}(e(u) - z) + \alpha \theta I + \text{Le}(\dot{u})\right):e(\dot{u})dxdt = \int_{Q_{\tau}} \dot{t} \cdot \dot{u}dxdt. \quad (7.1)$$

On the other hand, by using the definition of the subdifferential $\partial \Psi(z)$, we deduce from (3.9b) that

$$\int_{Q_{\tau}} \left(Mz:z + D_z H_1(z) + \theta D_z H_2(z) - \nu \Delta z\right)dxdt + \int_{Q_{\tau}} \Psi(z)dxdt = 0. \quad (7.2)$$

Adding (7.1) and (7.2), we obtain

$$\frac{1}{2} \int_{Q_{\tau}} \left(Mz:z + D_z H_1(z) - \theta D_z H_2(z) - \nu \Delta z\right)dxdt + \int_{Q_{\tau}} \Psi(z)dxdt = C^0 + \frac{1}{2} \int_{Q_{\tau}} \dot{t} \cdot \dot{u}dxdt. \quad (7.3)$$

Adding (7.1) and (7.2), we obtain

$$\frac{1}{2} \int_{Q_{\tau}} \left(Mz:z + D_z H_1(z) - \theta D_z H_2(z) - \nu \Delta z\right)dxdt + \int_{Q_{\tau}} \Psi(z)dxdt = C^0 + \frac{1}{2} \int_{Q_{\tau}} \dot{t} \cdot \dot{u}dxdt.$$
where \( C_0^{\alpha-z} \overset{\text{def}}{=} \frac{1}{T} \int_0^T E(e(u^0) - z^0); e(u^0) - z^0) \, dx + \frac{\gamma}{2} \| \nabla z^0 \|^2_{L^2(\Omega)} + \int_\Omega H_1(z^0) \, dx \). We integrate (3.6c) over \( Q_\tau \), by taking into account the boundary conditions (3.7), we find

\[
\int_\Omega \vartheta(\cdot, \tau) \, dx = \int_\Omega \vartheta^0 \, dx + \int_\Omega L(c(u)c(\hat{u})) \, dx + \int_\Omega M \ddot{z} \ddot{z} \, dx + \int_0^\tau \theta(\alpha r(e(\hat{u}))) \, dt + \int_{Q_\tau} \Phi(\ddot{z}) \, dx \, dt.
\]

We add this last equality to (7.3), and thanks to (3.11a), we obtain

\[
\frac{1}{2} \int_\Omega E(e(u^0; \tau) - z^0; e(u^0; \tau) - z^0) \, dx + \frac{\gamma}{2} \| \nabla z^0 \|^2_{L^2(\Omega)} + c_H \| z^0 \|^2_{L^2(\Omega)} + \int_\Omega \vartheta(\cdot, \tau) \, dx \leq C_0^{\alpha-z} + \| \vartheta^0 \|_{L^1(\Omega)} + c_H \| \Omega \| + \int_0^\tau \ell \cdot \dot{u} \, dx \, dt.
\]

Clearly there exists \( C_1 > 0 \), depending only on \( c, \nu \) and \( c, H_1 \) such that

\[
C_1 \| u(\cdot, \tau) \|^2_{H^1(\Omega)} + C_1 \| z(\cdot, \tau) \|^2_{H^1(\Omega)} + \int_\Omega \vartheta(\cdot, \tau) \, dx \leq C_0^{\alpha-z} + \| \vartheta^0 \|_{L^1(\Omega)} + c_H \| \Omega \| + \int_0^\tau \ell \cdot \dot{u} \, dx \, dt. \tag{7.4}
\]

Since \( \ell \in H^1(0, T; L^2(\Omega)) \), we may integrate by parts the last term of (7.4), we get

\[
\frac{1}{2} \int_\Omega E(e(u^0; \tau) - z^0; e(u^0; \tau) - z^0) \, dx + \frac{\gamma}{2} \| \nabla z^0 \|^2_{L^2(\Omega)} + \int_\Omega \vartheta(\cdot, \tau) \, dx \leq C_0^{\alpha-z} + \| \vartheta^0 \|_{L^1(\Omega)} + c_H \| \Omega \| + \int_0^\tau \ell \cdot \dot{u} \, dx \, dt.
\]

which allows us to conclude with Grönwall’s lemma since \( \vartheta(x, t) \geq 0 \) almost everywhere on \( Q_\tau \). \( \square \)

Note that (3.133) enable us also to obtain a global estimate for the temperature. More precisely, under the assumptions of Proposition (7.1), we have

\[
\| \vartheta \|_{L^\infty(0, \tau; L^2(\Omega))} \leq \left( \frac{\alpha}{\beta_1 C} \right)^{\frac{1}{2}}, \tag{7.5}
\]

for any solution \((u, z, \vartheta)\) of problem (3.1)-(3.3) defined on \([0, \tau] \) with \( \tau \in (0, T] \).

We assume that \( \beta_1 \geq 4 \). We define

\[
\bar{\theta} \overset{\text{def}}{=} T^{\frac{1}{4}} \| \theta \|_{L^\infty(0, \tau; L^2(\Omega))} \overset{\text{def}}{=} \left( \frac{\beta_1}{\alpha C} \right)^{\frac{1}{2}},
\]

and using the notations of Section 6, we define

\[
\bar{R}^f \overset{\text{def}}{=} R^f \left( \bar{R}^0, \bar{R}^\theta, \| \ell \|_{C^0([0, T]; L^2(\Omega))} \right), \quad \bar{R}^0 \overset{\text{def}}{=} C \left( T^{\frac{1}{4}} \bar{R}^f + \| \vartheta^0 \|_{L^2(\Omega)} \right), \quad \bar{R}^\theta \overset{\text{def}}{=} T^{\frac{1}{4}} \bar{R}^\infty + 1.
\]

Then, the results of Section 6 allow us to infer that there exists \( \tau \in (0, T] \) such that \( \phi_\tau \) admits a fixed point in \( B_{L^2(0, \tau; L^2(\Omega))}(0, \bar{R}^\theta) \). Let us define

\[
\bar{\tau} \overset{\text{def}}{=} \sup \{ \tau \in (0, T] : \phi_\tau \text{ admits a fixed point in } B_{L^2(0, \tau; L^2(\Omega))}(0, \bar{R}^\theta) \} \in (0, T].
\]

It is clear that problem (3.130)-(3.133) admits a global solution if and only if \( \bar{\tau} = T \). This identity is established below by a contradiction argument. To do so, we assume that \( \bar{\tau} \in (0, T) \) and we choose \( \epsilon > 0 \) such that \( \bar{\tau} - \epsilon \in (0, \bar{\tau}) \). By definition of \( \bar{\tau} \), there exists \( \tau \in (\bar{\tau} - \epsilon, \bar{\tau}] \) such that \( \phi_\tau \) admits a fixed point \( \bar{\vartheta} = \phi_\tau(\vartheta) \) in \( B_{L^2(0, T; L^2(\Omega))}(0, \bar{R}^\theta) \), i.e., the problem (3.10)-(3.13) admits a solution \((u, z, \vartheta)\) defined on \([0, \tau] \). We infer from the results of Section 7 that \( \| \vartheta \|_{L^\infty(0, \tau; L^2(\Omega))} \leq C \) and \( \| \vartheta(\cdot) \|_{L^2(0, \tau; L^2(\Omega))} \). Then, the definition of \( \bar{R}^\theta \) and the results of Section 8 imply that \( \bar{\vartheta} = \phi_\tau(\vartheta) \in L^\infty(0, \tau; L^2(\Omega)) \) with \( \| \vartheta \|_{L^\infty(0, \tau; L^2(\Omega))} \leq \bar{R}^\theta \).

Now let \( \bar{\tau} \in (0, T - \bar{\tau}) \) and \( \bar{R}^\theta \overset{\text{def}}{=} \left( \bar{R}^\theta \right)^{\hat{\theta}} \bar{\tau}(\bar{R}^\infty) \bar{\tau}^{\frac{\hat{\theta}}{4}} > 0 \). For any \( \hat{\vartheta} \in B_{L^2(0, T - \bar{\tau}; L^2(\Omega))}(0, \bar{R}^\theta) \), we define \( \bar{\vartheta}_{\text{ext}} \) as follows \( \bar{\vartheta}_{\text{ext}} \overset{\text{def}}{=} \vartheta \) on \([0, \tau] \), \( \bar{\vartheta}_{\text{ext}} \overset{\text{def}}{=} \hat{\vartheta} \) on \([\tau, \tau + \bar{\tau}] \) and \( \bar{\vartheta}_{\text{ext}} \overset{\text{def}}{=} 0 \) on \((\tau + \bar{\tau}, T] \). Clearly, we have

\[
\| \bar{\vartheta}_{\text{ext}} \|^2_{L^2(0, T; L^2(\Omega))} = \| \vartheta \|^2_{L^2(0, \tau; L^2(\Omega))} + \| \hat{\vartheta} \|^2_{L^2(\tau, \tau + \bar{\tau}; L^2(\Omega))} \leq \tau \bar{R}^\theta \bar{\tau} + (\bar{R}^\theta)^{\hat{\theta}} \leq (\bar{R}^\theta)^{\hat{\theta}}
\]
and the mapping \( \bar{\vartheta} \mapsto \bar{\vartheta}_{\text{ext}} \) is a contraction on \( L^q(\tau, \tau + \bar{\tau}; L^2(\Omega)) \). Let \( \bar{\vartheta} = \zeta(\bar{\vartheta}_{\text{ext}}) \). By definition of \( \zeta \), we have \( \bar{\vartheta} = \zeta(\vartheta) = \vartheta \) on \([0, \tau]\), \( \bar{\vartheta} = \zeta(\vartheta) \) on \( (\tau, \tau + \bar{\tau}) \) and \( \bar{\vartheta} = \zeta(0) = 0 \) on \((\tau + \bar{\tau}, T)\). Hence \( \bar{\vartheta} \in L^q(0, T; L^4(\Omega)) \) and

\[
\|\bar{\vartheta}\|_{L^q(0, T; L^4(\Omega))}^q \leq (\bar{R}^0)^q + \int_0^{\tau + \bar{\tau}} \|\zeta(\bar{\vartheta})\|_{L^4(\Omega)}^q \, dt \leq (\bar{R}^0)^q + \left( \frac{\beta_2}{\tilde{C}_{\text{ext}}} \right)^{\frac{q}{4}} \|\bar{\vartheta}\|_{L^q(\tau, \tau + \bar{\tau}; L^2(\Omega))} \|\bar{\vartheta}\|_{L^4(\Omega)}^\frac{q}{2},
\]

with \( \bar{R}^0 \equiv (\bar{R}^0)^q + \left( \frac{\beta_2}{\tilde{C}_{\text{ext}}} \right)^{\frac{q}{4}} (\bar{R}^0)^{\frac{q}{2}} \). By definition of \( \phi \), we get immediately that the restriction of \( \phi(\bar{\vartheta}_{\text{ext}}) \) on \([0, \tau]\) coincide with \( \phi_r(\vartheta) = \vartheta \) and we define \( \bar{\vartheta}_r(\bar{\vartheta}) \) as the restriction of \( \phi(\bar{\vartheta}_{\text{ext}}) \) to \((\tau, \tau + \bar{\tau})\). Furthermore, with the estimates of Section 6 we have \( \phi(\bar{\vartheta}_{\text{ext}}) \in L^\infty(0, T; L^2(\Omega)) \) and

\[
\|\phi(\bar{\vartheta}_{\text{ext}})\|_{L^\infty(0, T; L^2(\Omega))} \leq C(\bar{T}^{\frac{1}{4}} \bar{R}^0(0, \bar{R}^0), \|\phi\|_{C^0([0, T]; L^2(\Omega))} + \|\vartheta\|_{L^2(\Omega)}).
\]

It follows that there exists \( \bar{\tau} \in (0, T - \bar{\tau}) \), independent of \( \tau \), such that \( \bar{\vartheta}_r \) admits a fixed point \( \bar{\vartheta} \) in \( B_{L^q(\tau, \tau + \bar{\tau}; L^2(\Omega))}(0, \bar{R}^0) \). By construction of \( \bar{\vartheta}_r \), the restriction of \( \phi(\bar{\vartheta}_{\text{ext}}) \) to \([0, \tau + \bar{\tau}]\) is also a fixed point of \( \phi_{\tau + \bar{\tau}} \in B_{L^q(\tau, \tau + \bar{\tau}; L^2(\Omega))}(0, \bar{R}^0) \). But we may choose \( \epsilon \in (0, \bar{\tau}] \) such that \( \tau + \bar{\tau} > \bar{\tau} + \epsilon + \bar{\tau} > \bar{\tau} \), which gives a contradiction with the definition of \( \bar{\tau} \).

Hence we can conclude that \( \bar{\tau} = T \). Consequently, we deduce the following theorem:

**Theorem 7.2 (Global existence result)** Assume that (6.11), (6.12), (6.13), (6.16) and (6.17) hold. Assume moreover that \( \beta_1 \geq 4 \), \( u^0 \in V_0^q(\Omega) \), \( \vartheta^0 \in L^2(\Omega) \) such that (6.2) holds. Then the problem (3.8) admits a global solution \( (u, \vartheta, \vartheta) \) such that \( u \in W^{1, q}(0, T; V_0^q(\Omega)) \), \( z \in L^\infty(0, T; H^1(\Omega)) \), \( \Delta z \in L^2(0, T; L^2(\Omega)) \) and \( \vartheta \in C^0([0, T]; L^2(\Omega)) \) on \([0, T]\). Moreover \( \vartheta \) remains strictly positive and \( (u, z, \vartheta = \zeta(\vartheta)) \) is a solution of problem (3.1) on \([0, T]\).

**Remark 7.3** Let us assume furthermore that there exists \( \bar{C}_{H^2_z} > 0 \) and \( p_1 \in (0, 1) \) such that

\[
\forall z \in \mathbb{R}^{3x3}_{\text{dev}}, \quad |D_z H_2(z)| \leq \bar{C}_{H^2_z}(1 + |z|^{p_1}).
\]

Then we may obtain a global existence result for any \( \alpha > \max \left( 3, \frac{6}{5(1 - p_1)} \right) \). Indeed, we can establish a global a priori estimate for the enthalpy gradient by using a technique inspired by Boccardo and Gallouët [BocG90], i.e., by choosing the test-function \( \chi = 1 - \frac{1}{(1+\xi)^{d}} \) for some \( \xi > 0 \). By reproducing the same computations as in [Rou10, Prop. 4.2], we may obtain an estimate of \( \nabla \vartheta \) in \( L^r(0, \tau; L^L(\Omega)) \), independent of \( \tau \), for any \( r \in \left[ 1, \frac{d}{d + 2} \right] \) provided that \( \beta_1 > \frac{2d + 2}{d + 2(1 - p_1)} = \frac{2d + 2}{d + 2} \) since \( d = 3 \). Then we consider \( \alpha > 1 \) such that \( 2\mu \alpha \leq r \) and \( \frac{1}{\alpha} > \mu \left( \frac{1}{d} - \frac{1}{d + 1} \right) + 1 - \mu \) for some real number \( \mu \in (0, 1) \). Using Gagliardo-Nirenberg’s inequality, we may deduce that there exists \( C_{\text{GN}} > 0 \) such that

\[
\|\vartheta\|_{L^{2\alpha}(0, \tau; L^{L_0}(\Omega))} \leq C_{\text{GN}} \|\vartheta\|_{L^{(1 - \mu)2\alpha}(0, \tau; L^{L_0}(\Omega))} \int_0^\tau \left( \|\vartheta(t, \cdot)\|_{L^{L_0}(\Omega)} + \|\nabla \vartheta(t, \cdot)\|_{L^{L_0}(\Omega)} \right)^{2\alpha} \, dt,
\]

and reminding the estimate of \( \vartheta \) in \( L^\infty(0, \tau; L^2(\Omega)) \) obtained at Proposition 7.1 we infer that there exists a constant \( C > 0 \), depending only on the data, such that

\[
\|\vartheta\|_{L^{2\alpha}(0, \tau; L^{L_0}(\Omega))} \leq C \left( 1 + \|\nabla \vartheta\|_{L^{\infty}(0, \tau; L^{L_0}(\Omega))} \right).
\]

The three conditions

\[
1 \leq r < \frac{d + 2}{d + 1}, \quad 2\mu \alpha \leq r, \quad \frac{1}{\alpha} \geq \mu \left( \frac{1}{d} - \frac{1}{d + 1} \right) + 1 - \mu \quad \text{with} \quad 0 < \mu < 1,
\]

allow us to choose \( \mu = \frac{r d}{d + 2(1 - p_1)} \in (0, 1) \) and thus \( \alpha = \frac{d}{2\nu} = \frac{1}{2} + \frac{r(d + 1)}{2d} \in \left( \frac{4}{5}, \frac{4}{3} \right) \) (here \( d = 3 \)). It follows that, for any \( \beta_1 > \max(3, \frac{6}{5(1 - p_1)}) \), we will obtain a global estimate of \( \vartheta \) in \( L^{2\beta_1\alpha}(0, \tau; L^{\beta_1\alpha}(\Omega)) \) for any \( \alpha \in \left( \frac{4}{5}, \frac{4}{3} \right) \) and of \( \theta = \zeta(\vartheta) \) in \( L^{2\beta_1\alpha}(0, \tau; L^{\beta_1\alpha}(\Omega)) \). Thus, with \( \alpha \) such that \( \beta_1 \alpha > 4 \), we may obtain a global existence result by the same contradiction argument as in Section 7.

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