Carbon release by selective alloying of transition metal carbides

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Abstract
We have performed first principles density functional theory calculations on TiC alloyed on the Ti sublattice with 3d transition metals ranging from Sc to Zn. The theory is accompanied by experimental investigations, both as regards materials synthesis as well as characterization. Our results show that by dissolving a metal with a weak ability to form carbides, the stability of the alloy is lowered and a driving force for the release of carbon from the carbide is created. During thin film growth of a metal carbide this effect will favour the formation of a nanocomposite with carbide grains in a carbon matrix. The choice of alloying element as well as its concentration will affect the relative amount of carbon in the carbide and in the carbon matrix. This can be used to design the structures of nanocomposites and their physical and chemical properties. One example of applications is as low-friction coatings. Of the materials studied, we suggest the late 3d transition metals as the most promising elements for this phenomenon, at least when alloying with TiC.

(Some figures in this article are in colour only in the electronic version)

1. Introduction

Transition metal carbides (TMCs) have for a long time been the focus of extensive research, due to their interesting physical properties such as high melting points, hardness and conductivity. Many TMCs are also used as components in hard metal tools or as wear-resistant coatings on such tools [1]. A well known technique to deposit thin coatings of TMC is magnetron sputtering and one of the most studied TMC coating materials is TiC. With higher carbon contents the sputtered coatings usually form nanocomposites with nanocrystalline carbide grains in a carbon matrix or tissue phase, i.e. nc-TiC/a-C [2]. The properties of the TMC nanocomposite strongly depends on several factors such as the carbide grain size, thickness of the matrix phase and the relative distribution of carbide and matrix. Usually the growth process in sputtering is carried out far from equilibrium and it is therefore possible to alloy a carbide film with additional transition metal atoms by a substitutional solution on the Ti sublattice. The solubility of such metals in the sputtered coatings can be far higher than allowed at thermodynamical equilibrium, see e.g. [3–6], hence producing metastable phases. During annealing or under external pressures there will be a driving force for the metastable TMC coating to decompose into more stable phases. A schematic view of this concept is shown in figure 1, illustrating the decomposition process of a hypothetical metastable Ti\textsubscript{1−x}M\textsubscript{x}C\textsubscript{z} phase where the weak carbide forming metal M is substitutionally dissolved into the TiC structure. The thermodynamically most favourable decomposition route is always to remove the metal from the carbide structure by solid state diffusion and nucleation to form a metal-rich phase (left reaction route in figure 1). However, this requires substitutional diffusion of the metallic species which always has a much higher activation energy than interstitial diffusion. An alternative
and kinetically more favourable route is therefore interstitial diffusion of carbon towards the surface of the grains and a subsequent formation of free carbon (right reaction route in figure 1). Of course, this requires that it is energetically favourable for the ternary carbide to release some of its carbon. The exact values of the activation energies, $E_A$, in the $\text{Ti}_{1-x}M_xC$ phase will depend on the concentration of the metal M; however, literature data are known for the corresponding energies in pure TiC [7]. For this case, the activation energy for interstitial diffusion of carbon atoms has been determined to be 4.15 eV, which should be compared with substitutional diffusion of Ti with an activation energy of 7.65 eV [7]. We note that even though the activation energies for diffusion of the metal M and carbon in the ternary solid solutions are not known, there is no reason to assume that the activation energy for substitutional diffusion of the metal M will be lower than the activation energy for carbon diffusion in these systems.

The existence of two possible reaction routes will influence the growth of ternary, metastable carbides. Firstly, it suggests that a solid solution of a weak carbide forming metal into, e.g., TiC should favour the formation of free carbon between the grains and on the surface of the film. This has been observed experimentally by, e.g., Wilhelmsson et al [6, 8] for the case of alloying of Al into TiC. This effect can be used to form low-friction graphitic surfaces and design new types of low-friction materials [8]. Secondly, it also suggests that a process window should exist where at high temperatures a metal-rich phase is formed (left reaction route in figure 1), while at lower temperatures mainly carbon diffusion according to the right reaction route in figure 1 will be found.

In [8], we were able to demonstrate that metastable solid solutions of $\text{Ti}_{1-x}M_xC$ coatings could be deposited by magnetron sputtering and that, in agreement with theoretical predictions, they indeed formed surface graphitic layers upon both an annealing step and in tribological contacts. The surface layer of graphite reduced the friction coefficient to 0.05–0.1 using steel as a counter surface. An interesting observation is that coatings with a dissolved weak carbide forming metal are less hard but can exhibit similar wear properties compared to a pure TiC coating [6, 9, 10].

The concept of dissolving a weak carbide forming metal into a TMC to modify the structure and properties of a coating can be expanded to a large number of carbides and a wide range of dissolving metals. However, a systematic theoretical and experimental investigation of this effect has not yet been carried out. The stability of metal carbides has been studied by many authors, e.g. in [11–14]. The most important result within the present context is that the stability of the transition metal carbide is reduced with increasing number of d electrons. In fact, metals such as Fe, Ni and Cu form no thermodynamically stable carbides and it is therefore likely that different types of low-friction coatings can be designed by solid solutions of such weak carbide forming elements.

The aim of this study is to carry out a systematic investigation of the stability of a ternary solid solution of the metal M into TiC according to

$$\text{(1−x)Ti + (1−y)C + xM } \rightarrow \text{Ti}_{1−x}M_yC_{1−y}. \quad (1)$$

The stability of this solid solution has been investigated by calculations of the energy of formation with respect to Ti, M and C in their corresponding reference states. In this paper M has been restricted to the 3d transition metals series from Sc to Zn. Furthermore, the driving force for the out-diffusion of carbon under the formation of graphite has been calculated based on the energy difference between the phases in the reaction

$$\text{Ti}_{1-x}M_xC \rightarrow \text{Ti}_{1-x}M_yC_{1-y} + yC, \quad (2)$$

where C is in the form of graphite. This is equivalent to the evaluation of $\Delta E$ in figure 1. The results from the theoretical calculations have been compared with selected experimental data using Ni and Cu as alloying metals. As will be shown, the experimental results are in good agreement with the theoretical predictions.

The paper is arranged as follows: in section 2 we will discuss the necessary theoretical and experimental background, in section 3 we will discuss the theoretical results, in section 4 the experimental results are discussed and in section 5 we will summarize and draw conclusions.

2. Method

2.1. Theory

The TMCs that we focus on in this study all have the B1 (NaCl) structure, i.e. a face centred cubic metal lattice with C atoms occupying the octahedral interstitial positions. This is the ground state structure of TiC for a large compositional range when it comes to the C concentration [1, 15, 16]. For VC the B1 structure is also the ground state crystal structure but in a much narrower compositional range. For 3d transition metals from Cr to Zn there are essentially no chemically stable carbides, at least not in the B1 crystal structure [1]. However, since we are interested in the stability of alloyed TiC the B1 structure serves as a good reference crystal structure.

In order to model the ternary TMCs, the M atoms are substitutionally dissolved into TiC by substituting some of the Ti atoms on the metal lattice. This can be expanded to a large number of carbides and a wide range of dissolving metals. However, a systematic theoretical and experimental investigation of this effect has not yet been carried out. The stability of metal carbides has been studied by many authors, e.g. in [11–14]. The most important result within the present context is that the stability of the transition metal carbide is reduced with increasing number of d electrons. In fact, metals such as Fe, Ni and Cu form no thermodynamically stable carbides and it is therefore likely that different types of low-friction coatings can be designed by solid solutions of such weak carbide forming elements.

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In order to model the ternary TMCs, the M atoms are substitutionally dissolved into TiC by substituting some of the Ti atoms on the metal lattice. There are a number of issues related to the treatment of such a system. Assuming that the substitution on the metal lattice will leave the C
lattice unaffected, it is necessary to consider the effects of disorder on the metal lattice. Furthermore, since the TMC’s have relatively open structures, and due to the fact that the metals used in this study have different abilities to form carbides [1, 11–14], it is likely that the structure and stability of these carbides will depend on local environment effects such as local coordination and relaxations. Recently there has been a study on the importance of incorporating local environment effects in these types of systems [17]. Therefore, it is preferable to use a method that takes both disorder and local environment effects into account. Here we treat the problem of disorder within the supercell approach, where we have considered special quasi-random structures (SQS) that mimic the disorder of a random lattice [18]. The supercells that have been used in the calculations all contain 24 atomic positions on both the metal and the carbon sublattices. This means that in order to have 25% of the Ti atoms substituted by another metal atom, the supercell will contain 18 Ti and 6 M atoms. For 50% substitution, there will be 12 atoms of each species. For the carbon sublattice, the relation between C atoms and C vacancies works in a similar fashion to the metallic sublattice. Similar structures have been used by, for example, Alling et al [17]. The structures that have been used in this study are given in the appendix.

The calculations have been performed within density functional theory [19, 20] using the projector augmented wave method [21] as it is implemented in the Vienna ab initio simulation package [22, 23]. We have used the generalized gradient approximation for the description of the exchange and correlation functional [24]. The cut-off energy for the plane wave basis set was set to 600 eV in all calculations, and we have used the special k-points method of Monkhorst and Pack [25] and made sure that the calculations have been converged with respect to the number of k-points. Spin-polarized calculations have been performed for systems containing the metals Mn, Fe, Co and Ni; however, results related to the magnetic properties of these phases will be presented elsewhere. All calculations were fully relaxed in order to find the equilibrium volume and equilibrium atomic positions for each phase.

2.2. Experimental details

To verify the theoretical results, two sets of Ti–M–C films were deposited by non-reactive dc magnetron sputtering in an ultra-high vacuum chamber (base pressure below 1.0 × 10^{-7} Pa) using an Ar-discharge of 4.0 × 10^{-3} Pa. Elemental 2" targets of Ti (99.995% purity), M (>99.99% purity) and C (99.999% purity) were used. In the present study Ni and Cu were used as alloying elements. The films were deposited on substrates of Si(100). The films were characterized with x-ray photoelectron spectroscopy (XPS) and x-ray diffraction (XRD). XPS was performed using a Physical Electronics Quantum 2000 spectrometer and the samples were sputtered with 200 eV Ar\(^+\)ions prior to analysis to remove surface oxides. For all XPS analysis monochromatic Al K\(_\alpha\) radiation was used. The total composition of the samples was also determined using XPS data. XRD was performed using a Philips X’Pert MRD diffractometer, and grazing incidence (GI) scans were performed with an incident angle of 2°. Annealing was performed in a vacuum furnace at pressures below 7.0 × 10^{-5} Pa.

3. Theoretical results

We begin the presentation of our results by an analysis of the energy of formation, \(E_{\text{Form}}\), for the reaction given in equation (1), evaluated according to

\[
E_{\text{Form}} = E(\text{Ti}_{1-x}\text{M}_x\text{C}_{1-y}) - (1-x)E(\text{Ti}) - xE(M)
\]

(3)

where \(E(\text{Ti}_{1-x}\text{M}_x\text{C}_{1-y})\) is the total energy of the ternary carbide \(\text{Ti}_{1-x}\text{M}_x\text{C}_{1-y}\) and \(E(X)\) is the total energy of X, with X being either Ti, Sc or M. The energies of Ti, C and of the metal M are with respect to each element’s reference state in the bulk, see table 1. The results are presented in figure 2 for both stoichiometric \((y = 0)\) and sub-stoichiometric \((y = 0.25)\) phases. The calculations clearly show that the most stable carbide phase is pure TiC, and that by increasing the M concentration, \(x\), the formation energy increases. A positive formation energy corresponds to an unstable situation since the constituents will prefer to be separated in their corresponding single element phases. Note that there are only three binary carbide systems (TiC, ScC and VC) with a negative formation energy, signifying a good carbide forming ability. This is also in agreement with experimental data where only Ti, Sc and V form stable monocarbides with the B1 structure. As can be seen by comparing the upper and lower parts of figure 2, the sub-stoichiometric ScC\(_{0.75}\) is more stable than the stoichiometric ScC, which is also in agreement with experimental literature [26, 27]. Our results also show that alloying TiC with Sc and V on the metal sublattice results in negative formation energies for all values of the M content \(x\). Overall, the ternary carbides show a monotonic behaviour of the formation energy with respect to the alloying component.

An interesting parameter in figure 2 is the slope of the curve. The steeper the slope is, the more unfavourable a solid solution is, i.e. the less stable the ternary solution is. As can be

| Carbide   | \(y = 0\) | \(y = 0.25\) | Metal | Bulk reference |
|-----------|-----------|-------------|-------|----------------|
| ScC       | 4.68      | 4.72        | Sc    | hcp            |
| TiC       | 4.34      | 4.34        | Ti    | hcp            |
| VC        | 4.15      | 4.12        | V     | bcc            |
| CrC       | 4.07      | 4.00        | Cr    | bcc            |
| MnC       | 4.01      | 4.00        | Mn    | bcc            |
| FeC       | 3.99      | 3.94        | Fe    | bcc            |
| CoC       | 4.01      | 4.00        | Co    | bcc            |
| NiC       | 4.08      | 3.99        | Ni    | fcc            |
| CuC       | 4.24      | 4.16        | Cu    | fcc            |
| ZnC       | 4.40      | 4.36        | Zn    | fcc            |
|           |           |             |       | C              |

Carbide |
|---|
| Form | a (Å) |
|---|
| 0  | 4.68 |
| 0.25 | 4.72 |

Metals |
|---|
| Bulk reference |
|---|
| Sc | hcp |
| Ti | hcp |
| V  | bcc |
| Cr | bcc |
| Mn | bcc |
| Fe | bcc |
| Co | bcc |
| Ni | fcc |
| Cu | fcc |
| Zn | fcc |
Figure 2. The formation energy, $E_{\text{Form}}$, for the ternary solution Ti$_{1-x}$M$_x$C$_{1-y}$ according to equation (3) as a function of concentration, $x$, of the alloying element M for both full stoichiometric ($y = 0$, upper panel) and carbon deficient ($y = 0.25$, lower panel) phases. Energies are in units of eV/f.u., where f.u. stands for formula unit.

seen in figure 2, the variation in the slope is directly a function of the number of valence electrons in the system, which is in agreement with previous theoretical studies [11, 12]. Zn and Cu which form no carbides clearly give the most unfavourable solid solutions.

By comparing the stoichiometric and non-stoichiometric carbides we conclude that having a lower carbon concentration in the structure reduces the formation energy in general making the sub-stoichiometric ternary phases less unstable than the corresponding stoichiometric systems, as is readily shown in figure 2. This effect can be explained by the fact that in transition metal carbides the chemical bonding is known to be primarily due to strong metal–carbon bonds and a somewhat weaker metal–metal binding [28]. An in-mixing of late transition metals disrupts the strong covalent bonding between, e.g., Ti and C atoms in TiC, since the M–C bond is less favourable compared to the Ti–C bond, which increases the total energy of the system. If we now consider an increasing number of vacancies on the C lattice the importance of the metal–metal bond becomes more and more pronounced, simply because the number of C atoms to bind to is reduced. If the extreme case of 100% C vacancies is considered we are left with only metal–metal bonds in a face centred cubic lattice. So by alloying TiC with other transition metals, the strong metal–carbon bond becomes less favourable creating a tendency for the creation of C vacancies. The sub-stoichiometric phases on the other hand, while having a lower number of unfavourable metal–carbon bonds in the system, are also stabilized due to an increased metal–metal bonding in these phases. This change in behaviour of the bonding has also been shown to be reflected in the magnetic exchange interactions between Fe atoms in ternary Ti–Fe–C [29].

The equilibrium lattice constants of the ternary phases follow a Vegard’s law [30] behaviour, i.e. a linear variation, with respect to the alloying component, $x$. In table 1 we also show the evaluated lattice constants for the 3d binary carbide phases, for both stoichiometric ($y = 0$) and sub-stoichiometric ($y = 0.25$) phases. Generally, there is a slight decrease in the lattice parameter upon the introduction of vacancies, except for the case of ScC where there is a slight increase in the lattice parameter of the sub-stoichiometric phase compared to the stoichiometric case. An increased lattice parameter in sub-stoichiometric systems is a little surprising; however, the same behaviour can also be found for the B1 nitride phases of ZrN$_{1\text{--}y}$ and HfN$_{1\text{--}y}$ [1]. In the case of TiC, the lattice constant for the $y = 0.25$ phase is 1 mÅ shorter than for the $y = 0$ phase, which is in agreement with experimental
observations that the lattice parameter varies by less than 1% for $0 \leq y \leq 0.5$ [16, 31]. We also note that, for the binary systems examined here that exist in the B1 crystal structure, i.e. ScC$_{1-y}$, TiC$_{1-y}$ and VC$_{1-y}$, the evaluated lattice constants are in very good agreement with experimentally obtained lattice constants [1, 32].

We now turn our attention to the possibility of C leaving the system in the form of graphite as a response to the alloying with the metal M. The relevant property to consider from the theory is in this case the energy difference, $\Delta E$, which is illustrated in figure 1. This energy is given by the difference between the phases in the left and right parts of equation (2):

$$\Delta E = E(\text{Ti}_1 \rightarrow \text{M}_x \text{C}_{1-y}) + yE(C) - E(\text{Ti}_1 \rightarrow \text{M}_y \text{C}_x). \quad (4)$$

The results obtained by evaluating equation (4) for $y = 0.25$ are given in figure 3. Negative values signify that the carbide may lower its energy by releasing carbon in the form of graphite, i.e. that the right reaction route in figure 1 is a possible path for decomposition. Generally, graphite formation becomes more favourable by increasing the M concentration. However, in the case of M = V graphite is never formed, which is indicated by a slightly positive $\Delta E$ also for $x = 1$. This is consistent with the V–C phase diagram where a complete miscibility between TiC and VC is observed [33]. It is possible to conclude by looking at figure 3 that the ability for carbon to be released from the carbide increases with increasing M content. It should also be noted that the concentration at which carbon begins to be released from the carbide depends strongly on the M atom: for Sc and Cr $x \sim 0.5$ is needed, whereas for Ni, Cu and Zn $x \sim 0.3$ is needed. In fact, the curves in figure 3 are directly related to the stability curves in figure 2. The more unfavourable a solid solution is according to figure 2, the greater the driving force is to remove C from the structure and reduce the number of M–C bonds. Once more, this is a consequence of the decrease in the M–C bond strength with increasing number of valence electrons.

The results of figure 3 show that it is possible to cause a release of free C from TiC by dissolving several transition metal species on the metallic sublattice of TiC. This result follows previous theory and experiment, where only the substitution of Al was considered [8]. In the previous work, experimental studies demonstrated an increased rate of release of free C from Ti$_{1-x}$Al$_x$C$_2$ when Al was dissolved into the carbide and replacing some of the Ti on the metallic sublattice. This effect was also shown to influence the tribological properties of these carbide films, reducing the friction coefficient by more than 50% compared to the case of binary TiC [8]. Based on the results in figure 3 we suggest that a reduced friction coefficient should also be expected for solutions of other metallic species on the Ti sublattice of TiC. It is not unlikely that other transition metal carbides would show a similar pattern.

When regarding the electronic properties of the TMCs, it is known that the binary carbides show a more or less rigid band behaviour when changing the metal atom [11, 12]. In figure 4 we show calculated densities of states (DOSs) of Ti$_{1-x}$V$_x$C$_{1-y}$ as an illustration of how the electronic structure changes when varying the alloying component for both $y = 0$ and $y = 0.25$. For TiC (lowest curve in figure 4) the Fermi level is positioned in a region with a low DOS that separates bonding states at lower energies from anti-bonding states at higher energies, which have been formed by hybridization of p and d orbitals originating from the C and Ti atoms respectively [34]. All the bonding states are therefore occupied while the anti-bonding states at higher energies are empty. The replacement of Ti with other metals will effectively move the Fermi level and thereby the occupation of bonding and anti-bonding states depending on the alloying metal. This is the reason for the stability of TiC with respect to the other 3d TMCs; all bonding states are occupied while the anti-bonding states are not. In the ternary case, it is clear that the rigid band model can be used to explain the reduced stability of the ternary systems since the DOS is
not significantly different when changing the M content from $x = 0$ to 1, which is illustrated for the case of $Ti_{1-x}V_xC_{1-y}$ in figure 4. The reduced stabilities of the ternary phases are here due to the occupation of anti-bonding pd-hybridized states. Furthermore, the introduction of vacancies on the carbon lattice yields additional features in the DOS which arise from the lowered coordination of the metal atoms in the B1 crystal lattice for this situation. These features have a similar origin to the surface resonance states between bonding and anti-bonding states at the metal terminated (111)-surfaces on transition metal carbide surfaces [35].

4. Experimental results

The results in section 3 suggest that the addition of a weak carbide forming metal, substitutionally replacing Ti, should favour the formation of free carbon as graphite. Figure 5 shows diffractograms from as-deposited $Ti_{1-x}M_xC_{1-y}$ films with $M = Ni$ and Cu. As can be seen, diffraction peaks matching TiC appear in all diffractograms [36]. For the Ni containing series there is a clear shift of the diffraction peaks towards higher diffraction angles, signifying a decrease in lattice parameter from 4.38 to 4.32 Å as the Ni content increases. This is consistent with a substitutional solid solution of Ni in TiC; further details are published elsewhere [37, 38]. The Cu containing samples do not show the same decrease in lattice parameter. This may be explained by a lesser difference in atomic radii between Ti and Cu. Furthermore, the theoretical calculations predict that a solid solution carbide of Ti and Cu should be less stable than the solid solution carbide of Ti and Ni. This is also supported by XRD (see the upper right panel of figure 5) where the Cu containing samples exhibit a Cu phase, whereas the Ni containing samples synthesized under identical conditions show no Ni phase.

Experimentally there are two independent steps where the alloying of TiC with a weak carbide forming metal can lead to additional growth of free carbon. Firstly, during the film growth free carbon can be formed together with the carbide. In this case, the as-deposited films will exhibit a carbon phase together with the metastable carbide phase which has a solid
Annealing was performed under vacuum at 600 °C for 20 min. The ternary nc-(Ti, Ni)C/a-C sample clearly shows an increase in the amount of C phase, in contrast to the nc-TiC/a-C sample.

solution of the weak carbide forming metal. An increased driving force for graphite formation will then be seen as a larger relative number of C–C bonds compared to C–M and C–Ti bonds in, e.g., an XPS spectrum. Figure 5 also shows C 1s spectra of the same sample series as investigated using XRD. As can be seen, there is a pronounced peak at 281.8 eV which can be attributed to C–Ti bonds. This peak has a shoulder which can be attributed partly to C–M bonds and partly to a feature observed in nc-TiC/a-C which is related to interface effects [40, 41]. The peak at 284.2 eV is attributed to C–C bonds, and as can be seen the amount clearly increases as the M content increases. Since the total carbon content is approximately constant this means that the carbon content in the carbide phase decreases as a result of the alloying with the weak carbide forming metals Ni and Cu. Secondly, free carbon can form, due to the metastable structure of the coating, after film growth during, e.g., a high temperature annealing or as a result of external stress in, e.g., a tribological contact. Figure 6 shows the result of annealing experiments on nc-TiC/a-C and nc-(Ti, Ni)C/a-C thin films. It is clear that the amount of carbon phase is increased more in the case of the annealed ternary sample than in the case of the binary sample. It should be noted that no Ni phase could be observed by GI-XRD (not shown) after this annealing. Similar results have in previous studies also been attained for Ti–Al–C thin films [8].

5. Summary and conclusions

In summary, by using density functional theory calculations, we have identified several transition metal alloyed TiC systems where the introduction of weak carbide forming transition metals creates a tendency for the release of carbon from the structure. Energetically, carbon release is seen for all alloying metals that we have studied, except for alloying with V, but the release sets in for lower concentrations of dissolved metal in the cases of Fe, Co, Ni, Cu and Zn and for higher concentrations for Sc, Cr and Mn. The carbon release has been verified experimentally by XPS measurements on thin films for the case of TiC alloyed with Ni and Cu. We note that the out-diffusion of the metal M and a subsequent formation of metal-rich phases according to the left reaction route in figure 1 will always occur if the temperature is high enough to overcome the activation energy barriers for metal diffusion. However, based on the difference in experimentally evaluated activation energy barriers presented in section 1, there will be a process window where at lower temperatures the only possible route for decomposition is the release of carbon from the carbide phase.

The electronic structures of the ternary TMCs investigated in this study show a rigid band behaviour when varying the alloying component from lower to higher values. From the analysis of their electronic structures, we conclude that the reduced stability of the ternary TMCs arises from the occupation of anti-bonding states originating from hybridization between C p and metal d states. In the case of alloying with Sc, the reduced stability comes from the lesser occupation of the bonding states.

We conclude that for the purpose of creating a material for which the carbon release from the ternary carbide system can be utilized, e.g., a tribological contact, such as in the case of Ti1−xAlxC in [8], the late transition metals Ni, Cu and Zn would be the best candidates. We also note that due to the similarity of the chemical bonding of 3d TMCs to that in 4d and 5d TMCs, it is very likely that the behaviour presented here translates to the 4d and 5d TMCs. Furthermore, a number of papers on magnetron sputtering of ternary Ti–M–C films have clearly demonstrated that the addition of weak carbide forming elements such as Fe, Ni and Pt has a profound effect on the nanocomposite structure of as-deposited films [3–6, 37, 38, 42]. One of these effects is an increased amount of a matrix phase (free carbon) as demonstrated in figures 5 and 6. We suggest that this effect during the growth process itself is due to a decreased stability region (with regards to maximum carbon content) of TiC as it is alloyed with a weak carbide forming metal, since this reduces the number of unfavourable M–C bonds.

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Table A.1. Primitive lattice vectors, $a_i$, $i = 1, 2, 3$, and atomic coordinates for Ti, M and C used in the SQS supercells for $x = 0.25$ and 0.75 for both $y = 0.00$ and $y = 0.25$. The positions labelled with C(V) are empty in the case of $y = 0.25$. The short range order (SRO) parameters for the seven nearest shells are also given for each sublattice.

| Element | $x$   | $y$   | $z$   | Element | $x$   | $y$   | $z$   |
|---------|-------|-------|-------|---------|-------|-------|-------|
| $a_1$   | 0.500 | 0.500 | 0.000 | Ti      | 0.000 | 0.000 | 0.000 |
| $a_2$   | 0.000 | 0.375 | 0.375 | Ti      | 0.125 | 0.250 | 0.250 |
| $a_3$   | 0.250 | 0.000 | 0.250 | Ti      | 0.125 | 0.250 | 0.250 |

| Element | $x$   | $y$   | $z$   | Element | $x$   | $y$   | $z$   |
|---------|-------|-------|-------|---------|-------|-------|-------|
| Ti      | 0.000 | 0.000 | 0.125 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.125 | 0.250 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.250 | 0.000 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.125 | 0.250 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.250 | 0.125 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.125 | 0.250 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.250 | 0.250 | 0.125 | C       | 0.125 | 0.250 | 0.250 |

| Element | $x$   | $y$   | $z$   | Element | $x$   | $y$   | $z$   |
|---------|-------|-------|-------|---------|-------|-------|-------|
| Ti      | 0.125 | 0.250 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.250 | 0.000 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.125 | 0.250 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.250 | 0.125 | 0.250 | C       | 0.125 | 0.250 | 0.250 |
| Ti      | 0.125 | 0.125 | 0.000 | C(V)    | 0.125 | 0.125 | 0.125 |
| Ti      | 0.125 | 0.375 | 0.250 | C       | 0.125 | 0.375 | 0.250 |
| Ti      | 0.375 | 0.375 | 0.125 | C(V)    | 0.375 | 0.375 | 0.125 |
| Ti      | 0.250 | 0.500 | 0.250 | C       | 0.250 | 0.500 | 0.250 |

Table A.2. Primitive lattice vectors, $a_i$, $i = 1, 2, 3$, and atomic coordinates for Ti, M and C used in the SQS supercells for $x = 0.50$ for both $y = 0.00$ and $y = 0.25$. The positions labelled with C(V) are empty in the case of $y = 0.25$. The short range order (SRO) parameters for the seven nearest shells are also given for the metal lattice. The SRO parameters for the C lattice are given in table A.1.

| Element | $x$   | $y$   | $z$   | Element | $x$   | $y$   | $z$   |
|---------|-------|-------|-------|---------|-------|-------|-------|
| $a_1$   | 0.500 | 0.500 | 0.000 | Ti      | 0.000 | 0.000 | 0.000 |
| $a_2$   | 0.000 | 0.375 | 0.375 | Ti      | 0.375 | 0.375 | 0.000 |
| $a_3$   | 0.250 | 0.000 | 0.250 | Ti      | 0.250 | 0.250 | 0.250 |

| Element | $x$   | $y$   | $z$   | Element | $x$   | $y$   | $z$   |
|---------|-------|-------|-------|---------|-------|-------|-------|
| Ti      | 0.000 | 0.000 | 0.125 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.375 | 0.375 | 0.000 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.250 | 0.250 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.125 | 0.250 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.250 | 0.125 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.125 | 0.125 | 0.000 | C(V)    | 0.125 | 0.125 | 0.125 |
| Ti      | 0.125 | 0.375 | 0.250 | C       | 0.125 | 0.375 | 0.250 |
| Ti      | 0.375 | 0.375 | 0.125 | C(V)    | 0.375 | 0.375 | 0.125 |
| Ti      | 0.250 | 0.500 | 0.250 | C       | 0.250 | 0.500 | 0.250 |

| Element | $x$   | $y$   | $z$   | Element | $x$   | $y$   | $z$   |
|---------|-------|-------|-------|---------|-------|-------|-------|
| Ti      | 0.000 | 0.125 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.375 | 0.250 | 0.125 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.250 | 0.125 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| Ti      | 0.125 | 0.000 | 0.125 | C       | 0.125 | 0.125 | 0.125 |
| M       | 0.000 | 0.375 | 0.375 | C       | 0.125 | 0.125 | 0.125 |
| M       | 0.250 | 0.250 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
| M       | 0.125 | 0.125 | 0.000 | C(V)    | 0.125 | 0.125 | 0.125 |
| M       | 0.250 | 0.500 | 0.250 | C       | 0.125 | 0.125 | 0.125 |
Appendix

The distribution of Ti, M and C atoms in the special quasi-random structures that have been used in the calculations for \( x = 0.00, 0.25, 0.50, 0.75 \) and 1.00 for both \( y = 0.00 \) and 0.25 is shown in tables A.1 and A.2. For the cases of \( x = 0.25 \) and 0.75 the structures are identical except that the Ti and M atoms have been interchanged. The short range order (SRO) parameters between the Ti and M atoms on the metal lattice and between the C atoms and vacancies on the carbon lattice are also given in tables A.1 and A.2 for the seven nearest neighbour shells. Here the SRO parameters are defined such that the SRO is identical to zero for a given shell in the completely random system. For a recent thorough discussion on the definition of the SRO parameters given shell in the completely random system. A thorough discussion on the definition of the SRO parameters given shell in the completely random system. For a recent thorough discussion on the definition of the SRO parameters given shell in the completely random system.

The distribution of Ti, M and C atoms in the special quasi-random structures that have been used in the calculations for \( x = 0.00, 0.25, 0.50, 0.75 \) and 1.00 for both \( y = 0.00 \) and 0.25 is shown in tables A.1 and A.2. For the cases of \( x = 0.25 \) and 0.75 the structures are identical except that the Ti and M atoms have been interchanged. The short range order (SRO) parameters between the Ti and M atoms on the metal lattice and between the C atoms and vacancies on the carbon lattice are also given in tables A.1 and A.2 for the seven nearest neighbour shells. Here the SRO parameters are defined such that the SRO is identical to zero for a given shell in the completely random system. For a recent thorough discussion on the definition of the SRO parameters given shell in the completely random system. For a recent thorough discussion on the definition of the SRO parameters given shell in the completely random system.

Table A.2. (Continued.)

| Element | \( x \) | \( y \) | \( z \) | Element | \( x \) | \( y \) | \( z \) |
|---------|---------|---------|---------|---------|---------|---------|---------|
| M       | 0.125   | 0.125   | 0.250   | C(V)    | 0.375   | 0.500   | 0.500   |
| M       | 0.300   | 0.375   | 0.125   | C(V)    | 0.500   | 0.375   | 0.250   |
| M       | 0.375   | 0.625   | 0.250   | C(V)    | 0.375   | 0.500   | 0.250   |
| Shell   |         |         |         | Shell   |         |         |         |
| 1       | 0.000   | 0.000   | 0.000   | 2       | 0.000   | 0.000   | 0.000   |
| 3       | 0.000   | 0.000   | 0.000   | 4       | 0.000   | 0.000   | 0.000   |
| 5       | 0.000   | 0.000   | 0.000   |
| 6       | 0.000   | 0.000   | 0.000   |
| 7       | 0.000   | 0.000   | 0.000   |

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