Research Article

Analysis of the p-11B Fusion Scenario with Compensation of the Transfer of Kinetic Energy of Protons and Alpha Particles to the Gas Medium by the Electric Field

Mikhail L. Shmatov

Ioffe Institute, St. Petersburg 194021, Russia

Correspondence should be addressed to Mikhail L. Shmatov; m.shmatov@mail.ioffe.ru

Received 15 April 2022; Accepted 28 May 2022; Published 28 June 2022

Academic Editor: Dimitri Batani

Copyright © 2022 Mikhail L. Shmatov. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

The energy balance of the p-11B fusion scenario with compensation of the transfer of kinetic energy of protons and alpha particles to the gas medium by the electric field is considered. It is shown that such scenario cannot provide the use of p-11B fusion reaction for power production due to the very low ratio of the energy release of the fusion reaction to the energy necessary for compensation. The upper boundary of this ratio is about $2 \times 10^{-3}$.

1. Introduction

The influence of chain reactions on the rate $R_1$ of p-11B fusion reaction

$$p + ^{11}\text{B} \rightarrow 3\alpha + 8.7\text{ MeV}, \quad (1)$$

is discussed since 1973 [1–12]. One of the chain reactions consists of the scattering of at least one of the three alpha particles, generated by reaction (1), on proton(s) with acceleration of the proton(s) to kinetic energies, corresponding to a relatively high cross-section $\sigma_1$ for reaction (1) and the subsequent participation of the accelerated proton(s) in this reaction [1, 4, 6–12]. According to [1], at the temperature of 150–350 keV and the density of $10^{16}–10^{26}$ cm$^{-3}$, this chain reaction and other “nonthermal” effects result in an increase in $R_1$ on 5–15%. The type of particles with such densities was not mentioned [1], but this detail is not essential because in plasma under consideration, the densities of all particles are comparable [8]. According to [2–4, 7, 9, 10, 12], at least if special measures are taken, the increase in $R_1$ due to the chain reactions can be so high that it will provide the possibility of the use of reaction (1) for power production. The negative results of analysis of such assumptions from [2–4] are presented in [5, 6, 8, 11].

In 2020, Eliezer and Martinez-Val [9] and Eliezer et al. [10] proposed p-11B fusion scenarios with the influence of electric and magnetic fields on protons and alpha particles in the gas medium. The main idea of the proposal is that during some time periods, time-dependent electric field should compensate approximately for the transfer of kinetic energy $\epsilon_p$ of a proton with $\epsilon_p \approx \epsilon_p^\ast$, where $\epsilon_p^\ast$ is $\epsilon_p$ corresponding to the largest value of $\sigma_1$ for the collision of proton with the nucleus of $^{11}\text{B}$ in the rest, to the medium and for transfer of the kinetic energy of the alpha particle to the medium [9, 10]. This compensation should increase the probability of participation of the protons in reaction (1) and that of “useful” acceleration of protons due to the scattering of alpha particles on them. The magnetic field should provide the realization of these scenarios in reactors with acceptable sizes [9, 10]. Below, it is shown that in the scenario proposed in [10], the ratio $g$ of energy release of reaction (1) to the average value $\langle W_C \rangle$ of the energy spent for the initiation of one reaction (1) will be unacceptably low for power production.

2. The Upper Boundary of $g$

Eliezer et al. [10] analyzed the situations when reaction (1) occurs in gaseous $\text{H}_3^{11}\text{B}$ or other hydride of $^{11}\text{B}$ with a density of $10^{19}$ cm$^{-3}$ or of the order of $10^{19}$ cm$^{-3}$ and
temperature of about 1 eV or few eV. Ionization of this gas is supposed negligible [10]. Since a free molecule of H₃B does not exist and at the temperature above 700°C all hydrides of boron dissociate into boron and hydrogen [13], we will estimate the lowest boundary $W'_1$ of $\langle W'_1 \rangle$ in gas medium consisting of atoms of $^{11}$B with the density

$$n_{^{11}B} \approx 2.5 \times 10^{18} \text{ cm}^{-3},$$

and molecules of H₂ with the density

$$n_{\text{H}_2} \approx 3.75 \times 10^{18} \text{ cm}^{-3}.$$

At the conditions described in [10], the medium containing atoms of boron and molecules of hydrogen will also contain atoms of hydrogen and ions, but this is not essential for the analysis of the acceptability of attainable values of $g$ for power production. The ratio $n_{^{11}B}/n_{\text{H}_2}$ corresponds to the ratio of the numbers of nuclei of $^{11}$B and protons in the nonexisting free molecule of H₃$^{11}$B discussed in [10]. The choice of $n_{^{11}B}$ corresponds to an example presented on page 5 of Reference [10] and is mainly important for an estimate of the typical proton path $l_{\text{typ}} = 1/(\sigma_{^{11}B} n_{^{11}B})$, corresponding to one reaction (1). The estimate of $W'_1$ presented below yields that this parameter is independent of $n_{^{11}B}$.

In the situation under consideration, the change $\Delta P$ of $P$ on proton path $dx$ is given approximately by

$$\Delta P \approx \left[ eE - k_{^{11}B}^{\text{p}}(P) n_{^{11}B} - k_{\text{H}_2}^{\text{p}}(P) n_{\text{H}_2} \right] dx,$$

where $e$ is the proton charge, $E$ is the strength of the electric field, and $k_{^{11}B}^{\text{p}}$ and $k_{\text{H}_2}^{\text{p}}$ are the parameters describing the transfer of $\Delta P$ to molecules of hydrogen and atoms of boron, respectively. The parameter $k_{^{11}B}^{\text{p}}$ was calculated as

$$k_{^{11}B}^{\text{p}} = 2 A_{^{11}B} n_u S_{^{11}B},$$

where $A_{^{11}B}$ is the atomic mass of hydrogen, $n_u$ is the atomic mass unit, and $S_{^{11}B}$ is the stopping power of molecular hydrogen for proton. The parameter $k_{\text{H}_2}^{\text{p}}$ was calculated as

$$k_{\text{H}_2}^{\text{p}} = \frac{m_u}{2} \left( A_{\text{Be}} S_{\text{Be}}^{\text{p}} + A_{\text{C}} S_{\text{Cam}}^{\text{p}} \right),$$

where $A_{\text{Be}}$ is the atomic mass of beryllium, $S_{\text{Be}}^{\text{p}}$ is its stopping power for proton, $A_{\text{C}}$ is the atomic mass of carbon, and $S_{\text{Cam}}^{\text{p}}$ is the stopping power of amorphous carbon with the density of 2 g/cm³ for proton. The values of $S_{^{11}B}$, $S_{\text{Be}}^{\text{p}}$, and $S_{\text{Cam}}^{\text{p}}$ from [14] were used.

The parameter $k_{\text{Be}}^{\text{p}}$ was approximated by (6) due to the absence of data on the stopping power of boron for proton in [14]. This equation corresponds to the assumption that the product $P$ of the stopping power of the medium, consisting of atoms or molecules of one chemical element with atomic number $Z$, on the atomic mass of this element depends on $Z$ approximately linearly and, therefore,

$$P(Z) \approx [P(Z - \Delta Z) + P(Z + \Delta Z)]/2,$$

where $\Delta Z$ is a small natural number, for example, unity or two. In order to demonstrate that at least in some situations, the accuracy of (7) is rather high, let us compare $P(Z = 6, \varepsilon_p = 600 \text{keV}) = 3797 \text{MeV cm}^2 \text{ g}^{-1}$ and $P(Z = 6, \varepsilon_p = 700 \text{keV}) = 3440 \text{MeV cm}^2 \text{ g}^{-1}$, calculated using $S_{\text{Cam}}^{\text{p}}$ from [14], with the same parameters, calculated using (7) and $\Delta Z = 2$. Substituting $S_{\text{Be}}^{\text{p}}$ and the stopping power of molecular oxygen for proton from [14] into (7), we obtain $P(Z = 6, \varepsilon_p = 600 \text{keV}) = 3773 \text{MeV cm}^2 \text{ g}^{-1}$ and $P(Z = 6, \varepsilon_p = 700 \text{keV}) = 3424 \text{MeV cm}^2 \text{ g}^{-1}$. Thus, in these cases, the relative accuracy of (7) is better than 1%. This allows us to assume that at 600 keV $\leq \varepsilon_p \leq 700 \text{keV}$ (see below), the relative accuracy of (6) is of the order of 1% or even better.

According to [15], $\varepsilon_p \approx 646.2 \text{keV}$ and

$$\sigma_1(\varepsilon_p = \varepsilon_p^*) \approx 1.196 \text{b}.$$

Let us denote the value of $E$ corresponding to the condition $\Delta P/\Delta x = 0$, i.e., to the almost exact compensation of the transfer of kinetic energy of protons to the gas medium by the electric field, as $E_0$. This value depends on $\varepsilon_p$ (4). Equations (2)–(6) and (8)) yield that at $\varepsilon_p = \varepsilon_p^*$, $l_{\text{typ}} \approx 3.34 \times 10^5 \text{ cm}$, $E_0 \approx 24.9 \text{kV/cm}$, $e E_0 l_{\text{typ}} \approx 3424 \text{MeV cm}^2 \text{ g}^{-1}$, and $E_0 l_{\text{typ}} \approx 3424 \text{MeV cm}^2 \text{ g}^{-1}$. Thus, in these cases, the relative accuracy of (7) is better than 1%. This allows us to assume that at 600 keV $\leq \varepsilon_p \leq 700 \text{keV}$ (see below), the relative accuracy of (6) is of the order of 1% or even better.

According to [15], $\varepsilon_p \approx 646.2 \text{keV}$ and

$$\sigma_1(\varepsilon_p = \varepsilon_p^*) \approx 1.196 \text{b}.$$

Let us denote the value of $E$ corresponding to the condition $\Delta P/\Delta x = 0$, i.e., to the almost exact compensation of the transfer of kinetic energy of protons to the gas medium by the electric field, as $E_0$. This value depends on $\varepsilon_p$ (4). Equations (2)–(6) and (8)) yield that at $\varepsilon_p = \varepsilon_p^*$, $l_{\text{typ}} \approx 3.34 \times 10^5 \text{ cm}$, $E_0 \approx 24.9 \text{kV/cm}$, $e E_0 l_{\text{typ}} \approx 3424 \text{MeV cm}^2 \text{ g}^{-1}$, and $E_0 l_{\text{typ}} \approx 3424 \text{MeV cm}^2 \text{ g}^{-1}$. Thus, in these cases, the relative accuracy of (7) is better than 1%. This allows us to assume that at 600 keV $\leq \varepsilon_p \leq 700 \text{keV}$ (see below), the relative accuracy of (6) is of the order of 1% or even better.

According to [15], $\varepsilon_p \approx 646.2 \text{keV}$ and

$$\sigma_1(\varepsilon_p = \varepsilon_p^*) \approx 1.196 \text{b}.$$

Let us denote the value of $E$ corresponding to the condition $\Delta P/\Delta x = 0$, i.e., to the almost exact compensation of the transfer of kinetic energy of protons to the gas medium by the electric field, as $E_0$. This value depends on $\varepsilon_p$ (4). Equations (2)–(6) and (8)) yield that at $\varepsilon_p = \varepsilon_p^*$, $l_{\text{typ}} \approx 3.34 \times 10^5 \text{ cm}$, $E_0 \approx 24.9 \text{kV/cm}$, $e E_0 l_{\text{typ}} \approx 3424 \text{MeV cm}^2 \text{ g}^{-1}$, and $E_0 l_{\text{typ}} \approx 3424 \text{MeV cm}^2 \text{ g}^{-1}$. Thus, in these cases, the relative accuracy of (7) is better than 1%. This allows us to assume that at 600 keV $\leq \varepsilon_p \leq 700 \text{keV}$ (see below), the relative accuracy of (6) is of the order of 1% or even better.

According to [15], $\varepsilon_p \approx 646.2 \text{keV}$ and

$$\sigma_1(\varepsilon_p = \varepsilon_p^*) \approx 1.196 \text{b}.$$
when the product of the target gain on the driver efficiency $\eta_d$ exceeds ten. The target gain is the ratio of fusion energy release of one microexplosion to the energy delivered to the target for ignition of the microexplosion [20]. This parameter should exceed ten even if $\eta_d$ is close to unity and is an analog of the parameter $g$. Thus, $g$ of the order of $10^{-3}$ and less is not sufficient for power production involving conversion of fusion energy into thermal energy. Note that Weaver et al. [1] discussed briefly the potential feasibility of power production in the regime of subignition operation corresponding to $g < 1$. In any case, $g$ of the order of $10^{-3}$ and less seems to be too low even for this regime.

Note also that in the scenario proposed in [10], the acceleration of alpha particles, if it is not suppressed by the magnetic field, will not provide the effective acceleration of protons and, therefore, will serve mainly as a process increasing $\langle W_{\nu} \rangle$. This can be shown using equations, similar to (4)–(6), and the data from [10, 11, 14, 21] for the analysis of the motion of alpha particles and the transfer of their kinetic energy to protons. The compensation of deceleration of protons in the gas medium consisting mainly of atoms of $^{11}$B will also not provide sufficiently high values of $g$: at $n_{H2}/n_{11B2}$ = 0, the highest value of $1/(eE_d I_{typ})$ corresponds to $\epsilon_n \approx 656.6$ keV, $eE_d I_{typ} = 4.30$ GeV, and $(8.7$ MeV$/eE_0 I_{typ}) \approx 2.024 \times 10^{-3}$.

3. Conclusion

The scenario proposed in [10] cannot be used for effective power production due to the very low attainable $g$, the upper boundary of which is about $10^{-3}$. A decrease in $n_{H2}/n_{11B2}$ down to zero can result only in an approximately two-fold increase in the upper boundary of $g$. The real value of $g$ can be much less than its upper boundary.

Data Availability

The data used to support the findings of this study are included within the article.

Conflicts of Interest

The author declares that there are no conflicts of interest.

Acknowledgments

The author would like to thank the company HB11 Energy Pty Ltd. for the payment of APC for the publication of this manuscript.

References

[1] T. Weaver, G. Zimmerman, and L. Wood, Exotic CTR Fuels: Non-thermal Effects and Laser Fusion Applications,” Preprint UCRL-74938, Lawrence Livermore Laboratory, Livermore, CA, USA, 1973.

[2] V. S. Belyaev, V. P. Krainov, B. V. Zagreev, and A. P. Matafonov, “On the implementation of a chain nuclear reaction of thermonuclear fusion on the basis of the $p + ^{11}$B process,” Physics of Atomic Nuclei, vol. 78, no. 5, pp. 537–547, 2015.

[3] V. S. Belyaev, V. P. Krainov, A. P. Matafonov, and B. V. Zagreev, “The new possibility of the fusion $p + ^{11}$B chain reaction being induced by intense laser pulses,” Laser Physics Letters, vol. 12, no. 9, Article ID 096001, 2015.

[4] S. Eliezer, H. Hora, G. Korn, N. Nissim, and J. M. Martinez Val, “Avalanche proton-boron fusion based on elastic nuclear collisions,” Physics of Plasmas, vol. 23, no. 5, Article ID 050704, 2016.

[5] M. L. Shmatov, “Suppression of the chain nuclear fusion reaction based on the $p + ^{11}$B reaction because of the deceleration of alpha particles,” Physics of Atomic Nuclei, vol. 79, no. 5, pp. 666–670, 2016.

[6] M. L. Shmatov, “Comment on “Avalanche proton-boron fusion based on elastic nuclear collisions”,” Physics of Plasmas, vol. 23, no. 9, Article ID 094703, 2016.

[7] S. Eliezer, H. Hora, G. Korn, N. Nissim, and J. M. Martinez Val, “Response to “Comment on “Avalanche proton-boron fusion based on elastic nuclear collisions”,” Physics of Plasmas, vol. 23, no. 9, Article ID 094704, 2016.

[8] M. L. Shmatov, “Igniting a microexplosion by a microexplosion and some other controlled thermonuclear fusion scenarios with neutronless reactions,” Physics–Uspekhi, vol. 62, no. 1, pp. 70–81, 2019.

[9] S. Eliezer and J. M. Martinez-Val, “A novel fusion reactor with chain reactions for proton-boron11,” Laser and Particle Beams, vol. 38, no. 1, pp. 39–44, 2020.

[10] S. Eliezer, Y. Schweitzer, N. Nissim, and J. M. Martinez Val, “Mitigation of the stopping power effect on proton-boron11 nuclear fusion chain reactions,” Frontiers in Physics, vol. 8, Article ID 573694, 2020.

[11] F. Belloni, “On a fusion chain reaction via suprathermal ions in high-density $^{11}$B plasma,” Plasma Physics and Controlled Fusion, vol. 63, no. 5, Article ID 055020, 2021.

[12] J. Gruenwald, “On fusion chain reactions in $^{11}$B targets for laser driven aneutronic fusion,” Journal of Technological and Space Plasmas, vol. 2, no. 1, pp. 104–108, 2021.

[13] V. L. Vasilevskiy, Borovodorodki ['Hydrides of Boron'], in Bol’shaya Sovetskaya Encyclopedia [Big Soviet Encyclopedia, Sovetskaya Encyclopedia, Moscow, Russia, 1970].

[14] M. J. Berger, J. S. Coursey, M. A. Zuker, and J. Chang, “Stopping-power and range tables for electrons, protons and helium ions,” 2009, https://www.nist.gov/pml/data/star.

[15] W. M. Nevins and R. Swain, “The thermonuclear fusion rate coefficient for $p + ^{11}$B reactions,” Nuclear Fusion, vol. 40, no. 4, pp. 865–872, 2000.

[16] L. Nagy and L. Végh, “Ionization of molecular hydrogen by proton impact. I. Single ionization,” Physical Review A, vol. 46, no. 1, pp. 284–289, 1992.

[17] L. Nagy and L. Végh, “Ionization of molecular hydrogen by proton impact. II. Two-electron processes,” Physical Review A, vol. 46, no. 1, pp. 290–295, 1992.

[18] Y. P. Raizer, Fizika Gazovogo Razryada [Physics of Gas Discharge], Izd. Dom Intellekt, Dolgoprudny, Russia, 2009.

[19] E. Teller, “A future ICE (thermonuclear, that is!),” IEEE Spectrum, vol. 10, no. 1, pp. 60–64, 1973.

[20] V. P. Smirnov, V. I. Subbotin, and B. Y. Sharkov, Printispial’ Naya schema ITS” [“Principal scheme of ICF”], in Yadernyi Sintez S Inertnym Uderzhaniem [Nuclear Fusion with Inertial Confinement], Fizmatlit, Moscow, Russia, 2005.

[21] L. D. Landau and E. M. Lifshitz, Mechanics, Pergamon Press, Oxford, UK, 1962.