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DOI: https://doi.org/10.1007/JHEP06(2010)094
Calculation of the quark and gluon form factors to three loops in QCD

T. Gehrmann\textsuperscript{a}, E.W.N. Glover\textsuperscript{b}, T. Huber\textsuperscript{c}, N. Ikizlerli\textsuperscript{b}, C. Studerus\textsuperscript{a}

\textsuperscript{a} Institut für Theoretische Physik, Universität Zürich, Winterthurerstrasse 190, CH-8057 Zürich, Switzerland
\textsuperscript{b} Institute for Particle Physics Phenomenology, University of Durham, South Road, Durham DH1 3LE, England
\textsuperscript{c} Fachbereich 7, Universität Siegen, Walter-Flex-Strasse 3, D-57068 Siegen, Germany

ABSTRACT: We describe the calculation of the three-loop QCD corrections to quark and gluon form factors. The relevant three-loop Feynman diagrams are evaluated and the resulting three-loop Feynman integrals are reduced to a small set of known master integrals by using integration-by-parts relations. Our calculation confirms the recent results by Baikov et al. for the three-loop form factors. In addition, we derive the subleading $O(\epsilon)$ terms for the fermion-loop type contributions to the three-loop form factors which are required for the extraction of the fermionic contributions to the four-loop quark and gluon collinear anomalous dimensions. The finite parts of the form factors are used to determine the hard matching coefficients for the Drell-Yan process and inclusive Higgs-production in soft-collinear effective theory.

KEYWORDS: QCD, Multi-loop calculations.
1. Introduction

The form factors are basic vertex functions, and are as such fundamental ingredients for many precision calculations in QCD. They couple an external, colour-neutral off-shell current to a pair of partons: the quark form factor is the coupling of a virtual photon to a quark-antiquark pair, while the gluon form factor is the coupling of a Higgs boson to a pair of gluons through an effective Lagrangian. They appear as virtual higher-order corrections in coefficient functions for the inclusive Drell-Yan process [1, 2] and the inclusive Higgs production cross section [3–5]. In these observables, the infrared poles of the form factors cancel with infrared singularities from real radiation corrections. Consequently, it is possible to relate the coefficients of the infrared poles of the form factors to the coefficients of large logarithmic terms in the corresponding real radiation processes [6, 7]. A framework for combining the resummation of logarithmically enhanced terms at all orders with fixed-order results is provided in an effective field theory expansion [8] of QCD, which is systematized by soft-collinear effective theory [9]. In this context, the pole terms of the form factors yield the anomalous dimensions of the effective operators, while their finite terms determine the matching coefficients to a given order [10–12].
The form factors are actually the simplest QCD objects that display a non-trivial infrared pole structure. As such, their infrared pole coefficients can be used to extract fundamental constants: the cusp anomalous dimensions [13] which control the structure of soft divergences and the collinear quark and gluon anomalous dimensions. From the calculation [14, 15] of the pole terms of the three-loop form factors (and finite plus subleading terms in the two-loop and one-loop form factors [16–18]), these anomalous dimensions [14, 19, 20] are now known to three-loop order. An important observation is the agreement (up to an overall colour factor) of the cusp anomalous dimension for the quark and gluon, the so-called Casimir scaling [21]. Casimir scaling has been verified to three-loops [22, 23], but it is an open question whether it holds at four loops and beyond [24]. From non-perturbative arguments, the Casimir scaling is expected to break down at some loop order [25].

Based on the observation that infrared singularities of massless on-shell amplitudes in QCD are related to ultraviolet singularities of operators in soft-collinear effective field theory [13, 26], the pole structure of these amplitudes can be analyzed using operator renormalization. The singularity structure of arbitrary multi-leg massless QCD amplitudes is determined by an anomalous dimension matrix. The terms allowed in this anomalous dimension matrix are strongly constrained by relations between soft and collinear terms, from non-abelian exponentiation and from soft and collinear factorization. Independently, Becher and Neubert [20] and Gardi and Magnea [27] have proposed a remarkable all-loops conjecture that describes the pole structure of massless on-shell multi-loop multi-leg QCD amplitudes (generalizing earlier results at two [28] and three loops [29]) in terms of the cusp anomalous dimensions and the collinear anomalous dimensions. In this conjecture, the colour matrix structure of the soft anomalous dimension generated by soft gluons is simply a sum over two-body interactions between hard partons, and thus the matrix structure at any loop order is the same as at one loop. This result builds on the earlier work of Refs. [30, 31] which showed the colour matrix structure of the soft anomalous dimension at two loops is identical to that at one loop. There may be additional colour correlations at three loops or beyond, which cannot be excluded at present. However strong arguments for the absence of these terms are given in Refs. [20, 32]. If the all-order conjecture [20, 27] holds, the calculation of the pole parts of the form factors to a given loop order (and of the finite and subleading parts at fewer loops) would be sufficient to determine the infrared poles of all massless on-shell QCD amplitudes to this order.

The calculation of the three-loop form factors requires two principal ingredients: the algebraic reduction of all three-loop integrals appearing in the relevant Feynman diagrams to master integrals, and the analytical calculation of these master integrals. The reduction of integrals to master integrals exploits linear relations among different integrals, and is done based on a lexicographic ordering of the integrals (the Laporta algorithm [33]). Several dedicated computer-algebra implementations of the Laporta algorithm are available [33–36]. The reduction of the integrals relevant to the three-loop form factors is among the most challenging applications of the Laporta algorithm to date: due to the very large number of interconnected integrals to be reduced, the linear systems to be solved are often containing tens of thousand equations with a similar number of unknowns.

The master integrals in the three-loop form factors were identified already several years
ago [37]. Their analytical calculation proved to be a major computational challenge, which was completed only in several steps. The one-loop bubble insertions into two-loop vertex integrals as well as the two-loop bubble insertions into one-loop vertex integrals were derived using standard Feynman parameter integrals [37], while the genuine three-loop integrals required an extensive use of Mellin-Barnes integration techniques [38–40].

A first calculation of the three-loop form factors (based in part on numerical results for some of the expansion coefficients of the master integrals) was accomplished by Baikov et al. [41] in 2009. The analytical calculation of the last remaining master integrals was only completed recently [40]. It is the purpose of this paper to validate the three-loop form factor results of Ref. [40, 41] by an independent calculation, and to extend them in part to a higher order in the expansion in the dimensional regularization parameter \(\epsilon = 2 - d/2\). These further expansion terms will be needed for an extraction of the quark and gluon collinear anomalous dimensions from the single pole pieces of the four-loop form factors.

We define the quark and gluon form factors in Section 2, where we also discuss their UV-renomalization and summarize existing results at one- and two-loops. The reduction of the form factors to master integrals is described in Section 3, and the three-loop master integrals are discussed in Section 4. Explicit analytical expressions for them are collected in Appendix A. Our results for the three-loop form factors are presented in Section 5, and supplemented by Appendix B. The infrared structure of the QCD form factors up to four-loops is analyzed in Section 6. The three-loop hard matching coefficients for Drell-Yan and Higgs production in soft-collinear effective theory are determined from the form factors in Section 7. An outlook on future applications is contained in Section 8.

2. Quark and gluon form factors in perturbative QCD

The form factors are the basic vertex functions of an external off-shell current (with virtuality \(q^2 = s_{12}\)) coupling to a pair of partons with on-shell momenta \(p_1\) and \(p_2\). One distinguishes time-like (\(s_{12} > 0\), i.e. with partons both either in the initial or in the final state) and space-like (\(s_{12} < 0\), i.e. with one parton in the initial and one in the final state) configurations. The form factors are described in terms of scalar functions by contracting the respective vertex functions (evaluated in dimensional regularization with \(d = 4 - 2\epsilon\)) with projectors. For massless partons, the full vertex function is described with only a single form factor.

The quark form factor is obtained from the photon-quark-antiquark vertex \(\Gamma_{q\bar{q}}^\mu\) by

\[
\mathcal{F}^q = -\frac{1}{4(1-\epsilon)q^2} \text{Tr} \left( p_2^2 \Gamma_{\bar{q}q}^\mu \gamma_\mu \right),
\]

while the gluon form factor relates to the effective Higgs-gluon-gluon vertex \(\Gamma_{gg}^{\mu\nu}\) as

\[
\mathcal{F}^g = \frac{p_1 \cdot p_2 g_{\mu\nu} - p_{1,\mu} p_{2,\nu} - p_{1,\nu} p_{2,\mu}}{2(1-\epsilon)} \Gamma_{gg}^{\mu\nu}.
\]

The form factors are expanded in perturbative QCD in powers of the coupling constant, with each power corresponding to a virtual loop. We denote the unrenormalized form factors by \(\mathcal{F}^a\) and the renormalized form factors by \(F^a\) with \(a = q, g\).
At tree level, the Higgs boson does not couple either to the gluon or to massless quarks. In higher orders in perturbation theory, heavy quark loops introduce a coupling between the Higgs boson and gluons. In the limit of infinitely massive quarks, these loops give rise to an effective Lagrangian [42] mediating the coupling between the scalar Higgs field and the gluon field strength tensor:

\[ \mathcal{L}_{\text{int}} = -\frac{\lambda}{4} H F_{\mu\nu}^a F_{a,\mu\nu}. \]  

The coupling \( \lambda \) has inverse mass dimension. It can be computed by matching [43, 44] the effective theory to the full standard model cross sections [4].

Evaluation of the Feynman diagrams, contributing to the vertex functions at a given loop order yields the bare (unrenormalised) form factors,

\[ F_b^a(\alpha_s, s_{12}) = 1 + \sum_{n=1}^{\infty} \left( \frac{\alpha_s}{4\pi} \right)^n \left( \frac{-s_{12}}{\mu_0^2} \right)^{-n\epsilon} S_{\epsilon}^n F_n^a, \]

\[ F_b^b(\alpha_s, s_{12}) = \lambda_b \left( 1 + \sum_{n=1}^{\infty} \left( \frac{\alpha_s}{4\pi} \right)^n \left( \frac{-s_{12}}{\mu_0^2} \right)^{-n\epsilon} S_{\epsilon}^n F_n^b \right), \]

where \( \mu_0^2 \) is the mass parameter introduced in dimensional regularisation to maintain a dimensionless coupling in the bare Lagrangian density and where

\[ S_{\epsilon} = e^{-\gamma(4\pi)^\epsilon}, \]

with the Euler constant \( \gamma = 0.5772\ldots \)

The renormalization of the form factor is carried out by replacing the bare coupling \( \alpha_b \) with the renormalized coupling \( \alpha_s \equiv \alpha_s(\mu^2) \) evaluated at the renormalization scale \( \mu^2 \)

\[ \alpha_b \mu_0^{2\epsilon} = Z_{\alpha_s} \mu^{2\epsilon} \alpha_s(\mu^2). \]

For simplicity we set \( \mu^2 = |s_{12}| \) so that in the \( \overline{\text{MS}} \) scheme [45],

\[ Z_{\alpha_s} = S_{\epsilon}^{-1} \left[ 1 - \frac{\beta_0}{\epsilon} \left( \frac{\alpha_s}{4\pi} \right) + \left( \frac{\beta_0^2}{\epsilon^2} - 2 \beta_1 \right) \left( \frac{\alpha_s}{4\pi} \right)^2 \right. \]

\[ - \left. \left( \frac{\beta_0^3}{\epsilon^3} - \frac{7}{6} \frac{\beta_0 \beta_1}{\epsilon^2} + \frac{1}{3} \frac{\beta_2}{\epsilon} \right) \left( \frac{\alpha_s}{4\pi} \right)^3 + \mathcal{O}(\alpha_s^4) \right] , \]

where \( \beta_0, \beta_1 \) and \( \beta_2 \) are [46–48]

\[ \beta_0 = \frac{11 C_A}{3} - \frac{2 N_F}{3}, \]

\[ \beta_1 = \frac{34 C_A^2}{3} - \frac{10 C_A N_F}{3} - 2 C_F N_F, \]

\[ \beta_2 = \frac{2857 C_A^3}{54} + C_F^2 N_F - \frac{205 C_F C_A N_F}{18} - \frac{1415 C_A^2 N_F}{54} + \frac{11 C_F N_F^2}{9} + \frac{79 C_A N_F^2}{54}. \]

The renormalization relation for the effective coupling \( \lambda_b \) in the \( \overline{\text{MS}} \) scheme is given by,

\[ \lambda_b = \lambda \]

---
with
\[
Z_\lambda = 1 - \frac{\beta_0}{\epsilon} \left( \frac{\alpha_s}{4\pi} \right) + \left( \frac{\beta_0^2}{\epsilon^2} - \frac{\beta_1}{\epsilon} \right) \left( \frac{\alpha_s}{4\pi} \right)^2
- \left( \frac{\beta_0^3}{\epsilon^3} - \frac{2\beta_1\beta_0}{\epsilon^2} + \frac{\beta_2}{\epsilon} \right) \left( \frac{\alpha_s}{4\pi} \right)^3 + \mathcal{O}(\alpha_s^4).
\] (2.13)

The \( i \)-loop contribution to the unrenormalized coefficients is \( F^a_i \), while the renormalised coefficient is denoted by \( F^a_i \) where \( a = q, g \). If \( s_{12} \) is space-like, the form factors are real, while they acquire imaginary parts for time-like \( s_{12} \). These imaginary parts (and corresponding real parts) arise from the \( \epsilon \)-expansion of
\[
\Delta(s_{12}) = (-\text{sgn}(s_{12}) - i0)^{-\epsilon}
\] (2.14)
so that the renormalized form factors are given by,
\[
F^q(\alpha_s(\mu^2), s_{12}, \mu^2 = |s_{12}|) = 1 + \sum_{n=1}^{\infty} \left( \frac{\alpha_s(\mu^2)}{4\pi} \right)^n F^q_n,
\] (2.15)
\[
F^g(\alpha_s(\mu^2), s_{12}, \mu^2 = |s_{12}|) = \lambda \left( 1 + \sum_{n=1}^{\infty} \left( \frac{\alpha_s(\mu^2)}{4\pi} \right)^n F^g_n \right).
\] (2.16)

Up to three loops, the renormalized coefficients for the quark form factor (with \( \mu^2 = |s_{12}| \)) are then obtained as,
\[
F^q_1 = F^q_1 \Delta(s_{12}),
F^q_2 = F^q_2 (\Delta(s_{12}))^2 - \frac{\beta_0}{\epsilon} F^q_1 \Delta(s_{12}),
F^q_3 = F^q_3 (\Delta(s_{12}))^3 - \frac{2\beta_0}{\epsilon} F^q_2 (\Delta(s_{12}))^2 - \left( \frac{\beta_1}{2\epsilon} - \frac{\beta_0^2}{\epsilon^2} \right) F^q_1 \Delta(s_{12}),
\] (2.17)

while those for the gluon form factor are given by,
\[
F^g_1 = F^g_1 \Delta(s_{12}) - \frac{\beta_0}{\epsilon},
F^g_2 = F^g_2 (\Delta(s_{12}))^2 - \frac{2\beta_0}{\epsilon} F^g_1 \Delta(s_{12}) - \left( \frac{\beta_1}{\epsilon} - \frac{\beta_0^2}{\epsilon^2} \right),
F^g_3 = F^g_3 (\Delta(s_{12}))^3 - \frac{3\beta_0}{\epsilon} F^g_2 (\Delta(s_{12}))^2 - \left( \frac{3\beta_1}{2\epsilon} - \frac{3\beta_0^2}{\epsilon^2} \right) F^g_1 \Delta(s_{12}) - \left( \frac{\beta_2}{\epsilon} - \frac{2\beta_1\beta_0}{\epsilon^2} + \frac{\beta_3}{\epsilon^3} \right).
\] (2.18)

Unless explicitly stated otherwise, the renormalized form factors are given in the space-like case in the following sections.

The one-loop and two-loop form factors were computed in many places in the literature [14–18]. All-order expressions in terms of one-loop and two-loop master integrals are given in [18], and are summarized below.
2.1 Results at one-loop

Written in terms of the one-loop bubble integral, which is normalized to the factor

$$S_\Gamma = \frac{(4\pi)^\epsilon}{16\pi^2\Gamma(1-\epsilon)},$$  \hspace{1cm} (2.19)

the unrenormalised one-loop form factors are given by

$$\mathcal{F}_q^q/S_R = C_F B_{2,1} \left( \frac{4}{(D-4)} + D - 3 \right),$$  \hspace{1cm} (2.20)

$$\mathcal{F}_g^g/S_R = C_A B_{2,1} \left( \frac{4}{(D-4)} - \frac{4}{(D-2)} + 10 - D \right),$$  \hspace{1cm} (2.21)

where

$$S_R = \frac{16\pi^2 S_\Gamma}{S_\epsilon} = \frac{\exp(\epsilon\gamma)}{\Gamma(1-\epsilon)}.$$  \hspace{1cm} (2.22)

Eqs. (2.20) and (2.21) agree with eqs. (8) and (9) of ref. [18] respectively.

Inserting the expansion of the one-loop master integrals and keeping terms through to $O(\epsilon^5)$, we find that

$$\mathcal{F}_1^q = C_F \left[ -\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + (\zeta_2 - 8) + \epsilon \left( \frac{3 \zeta_2}{2} + \frac{14 \zeta_3}{3} - 16 \right) + \epsilon^2 \left( \frac{47 \zeta_2^2}{20} + 4 \zeta_2 + 7 \zeta_3 - 32 \right) \right.$$

$$+ \epsilon^3 \left( \frac{141 \zeta_2^3}{40} - \frac{7 \zeta_2 \zeta_3}{3} + 8 \zeta_2 + \frac{56 \zeta_3}{3} + \frac{62 \zeta_5}{5} - 64 \right)$$

$$+ \epsilon^4 \left( \frac{949 \zeta_2^3}{280} + \frac{47 \zeta_2^2}{5} - \frac{7 \zeta_2 \zeta_3}{2} - \frac{49 \zeta_3^2}{9} + 16 \zeta_2 + \frac{112 \zeta_3}{3} + \frac{93 \zeta_5}{5} - 128 \right)$$

$$+ \epsilon^5 \left( \frac{2847 \zeta_2^3}{560} + \frac{94 \zeta_2^2}{5} - \frac{329 \zeta_2 \zeta_3}{60} - \frac{28 \zeta_2 \zeta_5}{3} - \frac{31 \zeta_2 \zeta_5}{5} - \frac{49 \zeta_3^2}{6} \right.$$

$$+ 32 \zeta_2 + \frac{224 \zeta_3}{3} + \frac{248 \zeta_5}{5} + \frac{254 \zeta_7}{7} - 256 \right] \right),$$  \hspace{1cm} (2.23)

$$\mathcal{F}_1^g = C_A \left[ -\frac{2}{\epsilon^2} + \zeta_2 + \epsilon \left( \frac{14 \zeta_3}{3} - 2 \right) + \epsilon^2 \left( \frac{47 \zeta_2^2}{20} - 6 \right) \right.$$

$$+ \epsilon^3 \left( -\frac{7 \zeta_2 \zeta_3}{3} + \zeta_2 + \frac{62 \zeta_5}{5} - 14 \right)$$

$$+ \epsilon^4 \left( \frac{949 \zeta_2^3}{280} - \frac{49 \zeta_3^2}{9} + 3 \zeta_2 + \frac{14 \zeta_3}{3} - 30 \right)$$

$$+ \epsilon^5 \left( \frac{47 \zeta_2^3}{20} - \frac{329 \zeta_2 \zeta_3}{60} - \frac{31 \zeta_2 \zeta_5}{5} + 7 \zeta_2 + 14 \zeta_3 + \frac{254 \zeta_7}{7} - 62 \right) \right],$$  \hspace{1cm} (2.24)

where the gluon form factor agrees with eq. (7) of ref. [15] through to $O(\epsilon^4)$. Note that at each order in $\epsilon$, the terms of highest harmonic weight are the same for both quark and gluon form-factor. This is guaranteed by the equivalence of the coefficient of the leading pole in eqs. (2.20) and (2.21).
2.2 Results at two-loops

Written in terms of the two-loop master integrals (listed in the appendix), the unrenormalised two-loop gluon form factor is given by

\[
\frac{\mathcal{F}_g^2}{S_R^2} = C_F^2 \left[ B_{4,2} \left( \frac{16}{(D-4)^2} + \frac{8}{(D-4)} + D^2 - 6D + 17 \right) - C_{4,1} \left( \frac{7D^2}{8} - \frac{983D}{48} - \frac{565}{32(2D-7)} - \frac{20}{9(3D-8)} - \frac{28}{(D-4)} \right) \right.
\]

\[
\left. - \frac{40}{(D-4)^2} + \frac{10693}{288} \right] + B_{3,1} \left( \frac{27D^2}{8} - \frac{1293D}{16} + \frac{3955}{32(2D-7)} - \frac{17}{2(3D-3)} - \frac{476}{(D-4)} \right)
\]

\[
\left. - \frac{456}{(D-4)^2} - \frac{288}{(D-4)^3} + \frac{581}{32} \right]
\]

\[
- C_{6,2} \frac{D^3 - 20D^2 + 104D - 176}{8(2D-7)} \]

\[
+ C_F C_A \left[ - C_{4,1} \left( \frac{D^2}{16} + \frac{77D}{32} + \frac{565}{64(2D-7)} + \frac{12}{5(3D-8)} + \frac{23}{15(D-1)} \right) \right.
\]

\[
+ \frac{8}{3(D-4)} + \frac{16}{(D-4)^2} + \frac{163}{64} \right]
\]

\[
- B_{3,1} \left( \frac{75D^2}{16} - \frac{1837D}{32} + \frac{3955}{64(2D-7)} + \frac{3}{4(D-3)} - \frac{186}{(D-4)} \right)
\]

\[
- \frac{144}{(D-4)^2} - \frac{96}{(D-4)^3} + \frac{3845}{64} \right]
\]

\[
+ C_{6,2} \frac{D^3 - 20D^2 + 104D - 176}{16(2D-7)} \]

\[
+ C_F N_F \left[ - C_{4,1} \frac{(D-2)(3D^3 - 31D^2 + 110D - 128)}{(3D-8)(D-4)(D-1)} \right] \]

(2.25)

\[
\frac{\mathcal{F}_g^2}{S_R^2} = C_A^2 \left[ B_{4,2} \left( D^2 - 20D - \frac{48}{(D-2)} + \frac{32}{(D-4)} + \frac{16}{(D-2)^2} \right) \right.
\]

\[
+ \frac{16}{(D-4)^2} + \frac{100}{(D-4)^3} \right]
\]

\[
+ C_{4,1} \left( \frac{27D}{2} + \frac{119}{48(2D-5)} + \frac{75}{16(2D-7)} + \frac{10}{3(D-1)} + \frac{80}{(D-2)} \right)
\]

\[
+ \frac{103}{3(D-4)} - \frac{32}{(D-2)^2} + \frac{24}{(D-4)^2} - \frac{609}{8} \right]
\]

\[
+ B_{3,1} \left( 2AD + \frac{107}{144(2D-5)} + \frac{525}{16(2D-7)} + \frac{116}{9(D-1)} + \frac{96}{(D-2)} \right) \]
\[\frac{2}{(D-3)} - \frac{1175}{3(D-4)} - \frac{32}{(D-2)^2} - \frac{1388}{3(D-4)^2} - \frac{192}{(D-4)^3} - \frac{1955}{8} \]

\[+ C_{6,2} \frac{3(3D-8)(D-3)}{4(2D-5)(2D-7)} \]

\[+ C_{A N F} \left[ C_{4,1} \left( \frac{7D}{8} + \frac{119}{12(2D-5)} + \frac{35}{48(2D-7)} + \frac{20}{3(D-1)} - \frac{40}{3(D-2)} \right) \right. - \frac{2}{(D-4) - \frac{45}{16}} \]

\[-B_{3,1} \left( \frac{19D}{8} - \frac{107}{36(2D-5)} - \frac{245}{48(2D-7)} - \frac{232}{9(D-1) + 3(D-2)} \right) \]

\[-\frac{3}{2(D-3)} + \frac{8}{9(D-4)} - \frac{8}{(D-4)^2} - \frac{61}{16} \]

\[+ C_{6,2} \frac{(2D^3 - 25D^2 + 94D - 112)(D-4)}{8(D-2)(2D-5)(2D-7)} \]

\[+ C_{F N F} \left[ -C_{4,1} \frac{(46D^4 - 545D^3 + 2395D^2 - 4606D + 3248)(D-6)}{2(2D-7)(2D-5)(D-4)(D-2)} \right. \]

\[+ B_{3,1} \left( \frac{35D}{4} - \frac{107}{18(2D-5)} - \frac{245}{24(2D-7)} + \frac{8}{3(D-1) - \frac{1}{(D-3)}} \right) \]

\[-\frac{448}{9(D-4)} - \frac{112}{3(D-4)^2} - \frac{333}{8} \]

\[+ C_{6,2} \frac{(2D^3 - 25D^2 + 94D - 112)(D-4)}{4(D-2)(2D-5)(2D-7)} \]  \hspace{1cm} (2.26)

which, after re-expressing in terms of \(N\) and \(N_F\) agrees with eqs. (10) and (11) of ref. [18].

Inserting the expansion of the two-loop master integrals and keeping terms through to \(O(\epsilon^3)\), we find that

\[\mathcal{F}^2_2 = C_F^2 \left[ \frac{2}{\epsilon^3} + \frac{6}{\epsilon^3} - \frac{1}{\epsilon^2} \left( 2\zeta_2 - \frac{41}{2} \right) - \frac{1}{\epsilon} \left( \frac{64\zeta_3}{3} - \frac{221}{4} \right) \right. \]

\[- \left( 13\epsilon^2 - \frac{17\zeta_2}{2} + 58\zeta_3 - \frac{1151}{8} \right) \]

\[-\epsilon \left( \frac{171\zeta_2^2}{5} - \frac{112\zeta_2 \zeta_3}{3} - \frac{213\zeta_2}{4} + \frac{839\zeta_3}{3} + \frac{184\zeta_5}{5} - \frac{5741}{16} \right) \]

\[+ \epsilon^2 \left( \frac{223\zeta_2^3}{5} - \frac{3401\zeta_2^2}{20} + 54\zeta_2 \zeta_3 + \frac{2608\zeta_3^2}{9} + \frac{1839\zeta_2}{8} \right) \]

\[-\frac{6989\zeta_3}{6} + \frac{462\zeta_5}{5} + \frac{27911}{32} \]

\[+ \epsilon^3 \left( \frac{768\zeta_2^3}{7} + \frac{5488\zeta_2^2 \zeta_3}{15} - \frac{29157\zeta_2^2}{40} + \frac{757\zeta_2 \zeta_3}{3} + \frac{184\zeta_2 \zeta_5}{5} + \frac{2434\zeta_3^2}{3} \right) \]
\[ + C_F C_A \left[ -\frac{11}{6\epsilon^3} + \frac{1}{\epsilon^2} \left( \zeta_2 - \frac{83}{9} \right) - \frac{1}{\epsilon} \left( \frac{11\zeta_2}{6} - 13\zeta_3 + 4129/108 \right) + \left( \frac{44\zeta_2^2}{5} - \frac{119\zeta_2}{9} + \frac{467\zeta_3}{9} - 89173/648 \right) + \epsilon \left( \frac{1891\zeta_2}{60} - \frac{89\zeta_2\zeta_3}{3} - \frac{6505\zeta_2}{108} + \frac{6586\zeta_3}{27} + 51\zeta_5 - 1775893/3888 \right) - \epsilon^2 \left( \frac{809\zeta_2^3}{70} - \frac{2639\zeta_2^2}{18} + \frac{397\zeta_2\zeta_3}{9} + \frac{569\zeta_3^2}{3} + 146197\zeta_2/648 \right) - \frac{159949\zeta_3}{162} - \frac{3491\zeta_5}{15} + \frac{33912061}{23328} \right] + \epsilon^3 \left( \frac{3817\zeta_3^3}{140} - \frac{7103\zeta_2\zeta_3^2}{30} + \frac{638441\zeta_2^2\zeta_3}{1080} - \frac{4358\zeta_2\zeta_5}{27} - \frac{497\zeta_2\zeta_3}{5} - \frac{16439\zeta_2^3}{27} + \frac{2996725\zeta_2}{3888} + \frac{3709777\zeta_3}{972} + \frac{49786\zeta_5}{45} - \frac{372\zeta_7}{7} - \frac{63241291}{139968} \right), \tag{2.27} \]

which agrees through to \( O(\epsilon^2) \) with eq. (3.6) of ref. [14] and provides the next term in the expansion.

Similarly we find that the two-loop expansion of the gluon form factor is given by

\[ x_2^g = C_A^2 \left[ \frac{2}{\epsilon^3} - \frac{11}{6\epsilon^3} - \frac{1}{\epsilon^2} \left( \zeta_2 + \frac{67}{18} \right) + \frac{1}{\epsilon} \left( \frac{11\zeta_2}{2} - \frac{25\zeta_3}{3} + \frac{68}{27} \right) - \left( \frac{21\zeta_2^2}{5} - \frac{67\zeta_2}{6} - \frac{11\zeta_3}{9} - \frac{5861}{162} \right) - \epsilon \left( \frac{77\zeta_2^2}{60} - \frac{23\zeta_2\zeta_3}{3} - \frac{106\zeta_2}{9} + \frac{1139\zeta_3}{27} - \frac{71\zeta_5}{5} - \frac{158201}{972} \right) + \epsilon^2 \left( \frac{2313\zeta_2^3}{70} - \frac{1943\zeta_2^2}{60} - \frac{55\zeta_2\zeta_3}{3} + \frac{901\zeta_3^2}{9} + \frac{481\zeta_2}{54} - \frac{26218\zeta_3}{81} + \frac{341\zeta_5}{15} + \frac{3484193}{5832} \right) \right], \]
which agrees through to $O(\epsilon^2)$ with eq. (8) of ref. [15] and provides the next term in the expansion. Expressions for the renormalized one-loop and two-loop form factors, expanded to the appropriate order in $\epsilon$, can be found in [18].

3. Calculation of the three-loop form factors

To compute the three-loop quark and gluon form factors, we evaluate the relevant three-loop vertex functions within dimensional regularisation. At this loop order, there are 244 Feynman diagrams contributing to the quark form factor, and 1586 diagrams contributing to the gluon form factor. We generated these diagrams using QGRAF [49]. After contraction with the projectors (2.1)–(2.2), each diagram can be expressed as a linear combination of (typically hundreds of) scalar three-loop Feynman integrals. The three-loop integrals appearing in the form factors have up to nine different propagators. The integrands can depend on the three loop momenta, and the two on-shell external momenta, such that 12 different scalar products involving loop momenta can be formed. Consequently, not all scalar products can be cancelled against combinations of denominators, and we are left with irreducible scalar products in the numerator of the integrand. We denote the number of different propagators in an integral by $t$, the total number of propagators by $r$ and the
total number of irreducible scalar products by $s$. The topology of each integral is fixed by specifying the set of $t$ different propagators and subtopologies are obtained by removing one or more of the propagators.

Using relations between different integrals based on integration-by-parts (IBP) [50] and Lorentz invariance (LI) [51], one can express the large number of different integrals in terms of a small number of so-called master integrals. These identities yield large linear systems of equations, which are solved in an iterative manner using lexicographic ordering [33]. To carry out the reduction in a systematic manner, we introduce so-called auxiliary topologies. Each auxiliary topology is a set of 12 linearly independent propagators. Within the auxiliary topology, the integrand of a three-loop form factor integral with $(r, s, t)$ is expressed by $r$ propagators (with exactly $t$ different propagators) in the denominator, and $s$ propagators (with at most $12-t$ different propagators) in the numerator. All three-loop form factor integrals can be cast into one of three auxiliary topologies, which are listed in Table 1. The first auxiliary topology contains planar integrals only.

Three-loop integrals with $4 \leq t \leq 9$ and $t \leq r \leq 9$ appear in the form factors. These come with up to $s = 4$ irreducible scalar products for the quark form factor and up to $s = 5$ for the gluon form factor. For a fixed topology and given $(r, s, t)$, there are in total

$$N_{r,s,t} = \binom{r-1}{t-1} \binom{11-t+s}{s}$$

different integrals.

To obtain a reduction, one has to solve very large systems of equations. Already for $s \leq 4$, the system for a given auxiliary topology contains 900000 equations, and its solution is feasible only with dedicated computer algebra tools. For this reduction, we used the Mathematica-based package FIRE [35] and the C++ package Reduze [36], which was developed most recently by one of us.

| AuxTopo 1 | AuxTopo 2 | AuxTopo3 |
|-----------|-----------|----------|
| $k_1^2$   | $k_1^2$   | $k_1^2$  |
| $k_2^2$   | $k_2^2$   | $k_2^2$  |
| $k_3^2$   | $k_3^2$   | $k_3^2$  |
| $(k_1-k_2)^2$ | $(k_1-k_2)^2$ | $(k_1-k_2)^2$ |
| $(k_1-k_3)^2$ | $(k_1-k_3)^2$ | $(k_1-k_3)^2$ |
| $(k_2-k_3)^2$ | $(k_2-k_3)^2$ | $(k_1-k_2-k_3)^2$ |
| $(k_1-p_1)^2$ | $(k_1-k_3-p_2)^2$ | $(k_1-p_1)^2$ |
| $(k_1-p_1-p_2)^2$ | $(k_1-p_1-p_2)^2$ | $(k_1-p_1-p_2)^2$ |
| $(k_2-p_1)^2$ | $(k_2-p_1)^2$ | $(k_2-p_1)^2$ |
| $(k_2-p_1-p_2)^2$ | $(k_2-p_1-p_2)^2$ | $(k_2-p_1-p_2)^2$ |
| $(k_3-p_1)^2$ | $(k_3-p_1)^2$ | $(k_3-p_1)^2$ |
| $(k_3-p_1-p_2)^2$ | $(k_3-p_1-p_2)^2$ | $(k_3-p_1-p_2)^2$ |

Table 1: Propagators in the three different auxiliary topologies used to represent all three-loop form factor integrals.
With Reduze, the reduction and its performance are as follows. The topologies with more than 4 propagators are reduced after inserting the results of the sub-topologies into the system. With increasing $t$ the number of equations decrease as (in general) does the time taken to solve the system which is in the range of a few days to less than an hour with the program Reduze on a modern desktop computer. The total computing time for all the planar diagrams is more than 2 months. However, the parallelization of topologies with an equal number of propagators reduced the overall reduction time to a few weeks.

The three-loop form factors contain in total 22 master integrals, of which 14 are genuine three-loop vertex functions, 4 are three-loop propagator integrals and 4 are products of one-loop and two-loop integrals. They are described in detail in the following section.

4. Three-loop form factor master integrals

Our notation for the master integrals follows [37], and we distinguish three topological types of master integrals: genuine three-loop triangles ($A_{t,i}$-type), bubble integrals ($B_{t,i}$-type) and integrals that contain two-loop triangles ($C_{t,i}$-type). In this notation, the index $t$ denotes the number of propagators, and $i$ is simply enumerating the topologically different integrals with the same number of propagators.

The one-loop and two-loop master integrals appearing in the form factors at these loop orders are displayed in Figure 1. Their expansions to finite order have been known for a long time, all-orders expressions were derived in [18], they can for example be expanded using HypExp [52]. $B_{t,i}$-type and $C_{t,i}$-type three-loop integrals are listed in Figure 2. The $B_{t,i}$-type integrals were computed to finite order in [50, 53], and supplemented by the higher order terms in [54]. Finally, the genuine three-loop vertex integrals are shown in Figure 3, their expansions to finite order were derived in [37–40].

The calculation of the nine-liner three-loop integrals was the last missing ingredient to the form factor calculation for a long time. The full result for $A_{9,1}$ and most of the pole parts of $A_{9,2}$ and $A_{9,4}$ were computed analytically in [39]. Analytical expressions for the

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**Figure 1:** One and two-loop master integrals appearing in the quark and gluon form factors.
remaining pieces of the latter two integrals were subsequently obtained in [40]. In [39], it was pointed out that for each of these three integrals one can find an integral from the same topology with an irreducible scalar product, which has homogeneous transcendentality. These integrals were named $A_{9,1n}$, $A_{9,2n}$, and $A_{9,4n}$, and are defined in [39]. Compared to [39] we increased the numerical precision of the remaining coefficients, both for $A_{9,2}$ and $A_{9,4}$, by means of conventional packages like MB.m [55]. We reproduce thirteen significant digits of the analytic result of [40] in the case of $A_{9,2}$, and fourteen in the case of $A_{9,4}$. When converted to $A_{9,2n}$ and $A_{9,4n}$ in order to match onto a homogeneous basis of transcendental constants, we find a unique solution using PSLQ [56] for the simple pole of $A_{9,4}$, but we do not (yet) find a solution for the finite pieces of $A_{9,2}$ or $A_{9,4}$.

An analytic result for $A_{9,2}$ and $A_{9,4}$, derived by purely analytic steps and without fitting rational coefficients to numerical values, is still a desirable task, and remains to be investigated in the future. This goal is definitely within reach in the case of $A_{9,4}$, whereas the situation is less clear for $A_{9,2}$.

Expansions of all master integrals to the order in $\epsilon$ where transcendentality six first appears are listed in the Appendix.

5. Three-loop form factors

The unrenormalised three-loop form factors can be decomposed into different colour struc-
Figure 3: Three-point integrals listed in Refs. [37–39].

\[ \frac{F_3}{S_R} = C_F^3 X_C^q + C_F^2 C_A X_C^q C_A + C_F C_A^2 X_C^q C_A^2 + C_F N_F X_C^q N_F \]
\[ + C_F C_A N_F X_C^q C_A N_F + C_F N_F^2 X_C^q N_F^2 + C_F N_{F,V} \left( \frac{N^2 - 4}{N} \right) X_C^q N_{F,V} (5.1) \]
and

\[
\mathcal{F}_3^q / S_R^3 = C_A^3 X_{C_A}^q + C_A^2 N_F X_{C_A N_F}^q + C_A C_F N_F X_{C_A C_F N_F}^q + C_F^2 N_F X_{C_F N_F}^q + C_A N_F^2 X_{C_F N_F}^q + C_F N_F^2 X_{C_F N_F}^2 ,
\]

where the last term in the quark form factor is generated by graphs where the virtual gauge boson does not couple directly to the final-state quarks. This contribution is denoted by \(N_{F,V}\) and is proportional to the charge weighted sum of the quark flavours. In the case of purely electromagnetic interactions, we find,

\[
N_{F,\gamma} = \sum q e_q.
\]

The coefficient of each colour structure is a linear combination of master integrals, resulting from the reduction of the integrals appearing in the Feynman diagrams. All coefficients are listed in Appendix B.

Inserting the expansion of the three-loop master integrals and keeping terms through to \(O(\epsilon^0)\), we find that the three-loop coefficients are given by

\[
\mathcal{F}_3^q = C_F^3 \left[ -\frac{4}{3 \epsilon^6} + \frac{6}{\epsilon^5} + \frac{1}{\epsilon^4} (2 \zeta_2 - 25) + \frac{1}{\epsilon^3} \left( -3 \zeta_2 + \frac{100 \zeta_3}{3} - 83 \right) + \frac{1}{\epsilon^2} \left( \frac{213 \zeta_2}{10} - \frac{77 \zeta_2}{2} + 138 \zeta_3 - \frac{515}{2} \right) + \frac{1}{\epsilon} \left( \frac{1461 \zeta_2}{20} - \frac{214 \zeta_2 \zeta_3}{3} - \frac{467 \zeta_2}{2} + \frac{2119 \zeta_3}{3} + \frac{644 \zeta_5}{5} - \frac{9073}{12} \right) + \left( \frac{53675}{24} - \frac{13001 \zeta_2}{12} + \frac{12743 \zeta_2^2}{40} - \frac{9095 \zeta_3^3}{252} + 2669 \zeta_3 + 61 \zeta_3 \zeta_2 \right) - \frac{1826 \zeta_3^2}{3} + \frac{4238 \zeta_5}{5} \right] + C_F^2 C_A \left[ -\frac{11}{3 \epsilon^6} + \frac{1}{\epsilon^5} \left( -2 \zeta_2 + \frac{431}{18} \right) + \frac{1}{\epsilon^4} \left( -\frac{7 \zeta_2}{6} - 26 \zeta_3 + \frac{6415}{54} \right) + \frac{1}{\epsilon^3} \left( -\frac{83 \zeta_2}{5} + \frac{1487 \zeta_2}{36} - 210 \zeta_3 + \frac{79277}{162} \right) + \frac{1}{\epsilon^2} \left( -\frac{9839 \zeta_2}{72} + \frac{215 \zeta_2 \zeta_3}{3} + \frac{38623 \zeta_2}{108} - \frac{6703 \zeta_3}{6} - 142 \zeta_5 + \frac{1773839}{972} \right) + \left( \frac{37684115}{5832} + \frac{664323 \zeta_2}{324} - \frac{1265467 \zeta_2^2}{2160} - \frac{18619 \zeta_3^3}{1260} - \frac{96715 \zeta_3}{18} + \frac{46 \zeta_2 \zeta_3}{9} + \frac{161 \zeta_3^2}{3} - \frac{46594 \zeta_5}{45} \right) \right] + C_F C_A \left[ \frac{242}{81 \epsilon^4} + \frac{1}{\epsilon^3} \left( \frac{88 \zeta_2}{27} - \frac{6521}{243} \right) + \frac{1}{\epsilon^2} \left( -\frac{88 \zeta_2^2}{45} - \frac{553 \zeta_2}{81} + \frac{1672 \zeta_3}{27} - \frac{40289}{423} \right) + \frac{1}{\epsilon} \left( \frac{802 \zeta_2}{15} - \frac{88 \zeta_2 \zeta_3}{9} + \frac{68497 \zeta_2}{486} + \frac{12106 \zeta_3}{27} - \frac{136 \zeta_5}{3} - \frac{1870564}{2187} \right) \right] - 15 -
The pole contributions of $F_{q}^{3}$ are given in eq. (3.7) of ref. [14] while the finite parts of the $N_{F}^{2}$, $C_{A}N_{F}$ and $C_{F}N_{F}$ contributions are given in eq. (6) of ref. [15]. The finite $N_{F,V}$ contribution can be obtained from the $\delta(1-x)$ contribution to the $d_{abc}d_{abc}$ colour factor in eq. (6.6) of ref. [57]. The remaining finite contributions are given in eqs. (8) and (9) of ref. [41].

Similarly, the expansion of the gluon form factor at three-loops is given by

\[
F_{g}^{3} = C_{A}^{3} \left[ -\frac{4}{3e^{6}} + \frac{11}{3e^{6} + \frac{361}{81e^{4}}} + \frac{1}{e^{2}} \left( -\frac{517\zeta_{2}}{54} + \frac{22\zeta_{3}}{3} - \frac{3506}{243} \right) \\
+ \frac{1}{e^{2}} \left( \frac{247\zeta_{2}^{2}}{90} + \frac{481\zeta_{2}}{162} - \frac{209\zeta_{3}}{27} - \frac{17741}{243} \right) \\
+ \frac{1}{e} \left( \frac{3751\zeta_{2}^{2}}{360} - \frac{85\zeta_{2}\zeta_{3}}{9} + \frac{20329\zeta_{2}}{243} + \frac{241\zeta_{3}}{9} - \frac{878\zeta_{5}}{15} - \frac{145219}{2187} \right) \\
+ \left( \frac{14474131}{13122} + \frac{307057\zeta_{2}}{1458} + \frac{8459\zeta_{2}^{2}}{1080} - \frac{22523\zeta_{2}^{3}}{270} \right) \right].
\]
\[ \begin{aligned} + C_A N_F & \left[ - \frac{68590 \zeta_3}{243} + \frac{77 \zeta_2 \zeta_3}{18} - \frac{1766 \zeta_3^2}{9} + \frac{20911 \zeta_5}{45} \right] \\
+ C_A^2 N_F & \left[ - \frac{2}{3 \epsilon^2} - \frac{2}{81 \epsilon^4} + \frac{1}{\epsilon^3} \left( \frac{47 \zeta_2}{27} + \frac{1534}{243} \right) + \frac{1}{\epsilon^2} \left( - \frac{425 \zeta_2}{81} + \frac{518 \zeta_3}{27} + \frac{4280}{243} \right) + \frac{1}{\epsilon} \left( \frac{2453 \zeta_2^2}{180} - \frac{7561 \zeta_2}{243} + \frac{1022 \zeta_3}{81} - \frac{92449}{2187} \right) + \left( \frac{437 \zeta_2^2}{60} - \frac{439 \zeta_2 \zeta_3}{729} - \frac{37868 \zeta_2}{27} - \frac{754 \zeta_3}{27} + \frac{3238 \zeta_5}{45} - \frac{10021 \zeta_3 \zeta_5}{13122} \right) \right] \\
+ C_A C_F N_F & \left[ \frac{20}{9 \epsilon^3} + \frac{1}{\epsilon^2} \left( - \frac{160 \zeta_3}{9} + \frac{526}{27} \right) + \frac{1}{\epsilon} \left( - \frac{176 \zeta_3^2}{15} - \frac{22 \zeta_2}{3} - \frac{224 \zeta_3}{27} + \frac{2783}{81} \right) + \left( - \frac{16 \zeta_2^2}{5} + \frac{48 \zeta_2 \zeta_3}{3} - \frac{41 \zeta_2}{81} + \frac{11792 \zeta_3}{27} + \frac{32 \zeta_5}{9} - \frac{155629}{486} \right) \right] \\
+ C_F N_F & \left[ \frac{2}{3 \epsilon} + \frac{1}{\epsilon^2} \left( \frac{296 \zeta_3}{3} - 160 \zeta_3 + \frac{304}{9} \right) \right] \\
+ C_A N_F^2 & \left[ - \frac{8}{81 \epsilon^4} + \frac{80}{243 \epsilon^3} + \frac{1}{\epsilon^2} \left( \frac{20 \zeta_2}{27} + \frac{8}{9} \right) + \frac{1}{\epsilon} \left( \frac{200 \zeta_2}{81} + \frac{664 \zeta_3}{81} + \frac{34097}{2187} \right) + \left( \frac{797 \zeta_2^2}{135} + \frac{76 \zeta_2}{27} + \frac{11824 \zeta_3}{243} + \frac{1479109}{13122} \right) \right] \\
+ C_F N_F^2 & \left[ \frac{8}{9 \epsilon^2} + \frac{1}{\epsilon} \left( - \frac{32 \zeta_3}{3} + \frac{424}{27} \right) + \left( - \frac{112 \zeta_2^2}{15} - \frac{16 \zeta_2}{3} - \frac{704 \zeta_3}{9} + \frac{10562}{81} \right) \right]. \tag{5.5} \end{aligned} \]

The divergent parts agree with eq. (8) of ref. [15] while the finite contributions agree with eq. (10) of ref. [41].

Using our knowledge of the three-loop form factors, we can also write down the \( \mathcal{O}(\epsilon) \) contributions to the \( N_F \) parts of the quark and gluon form factors. For the quark form-factor we find that,

\[ \begin{aligned} \mathcal{F}_3^{\gamma |N_F} &= C_F N_F^2 \epsilon \left( - \frac{2913928}{6561} + \frac{2248}{135} \zeta_5 + \frac{2108}{27} \zeta_3 - \frac{24950}{243} \zeta_2 - \frac{68}{9} \zeta_2 \zeta_3 - \frac{3901}{810} \zeta_2^2 \right) \\
+ C_F C_A N_F \epsilon \left( \frac{24570881}{4374} - \frac{28156}{45} \zeta_5 - \frac{2418896}{729} \zeta_3 + \frac{10816}{27} \zeta_3^2 + \frac{7137385}{4374} \zeta_2 \right) + \frac{2674}{27} \zeta_2 \zeta_3 - \frac{352559}{1620} \zeta_2^2 + \frac{17324}{945} \zeta_2^3 \right) \\
+ C_F^2 N_F \epsilon \left( - \frac{50187205}{17496} + \frac{5863}{135} \zeta_5 + \frac{929587}{243} \zeta_3 - \frac{5771}{9} \zeta_3^2 - \frac{1263505}{972} \zeta_2 \right) - \frac{8515}{54} \zeta_2 \zeta_3 + \frac{821749}{3240} \zeta_2^2 - \frac{875381}{7560} \zeta_2^3 \right) \end{aligned} \]
\[ + C_F N_{F,V} \left( \frac{N^2 - 4}{N} \right) \varepsilon \left( \frac{170}{3} + \frac{752}{9} \zeta_5 + \frac{94}{9} \zeta_3 - \frac{344}{3} \zeta_5^2 + \frac{260}{3} \zeta_2 \right) \\
+ 30 \zeta_2 \zeta_3 - \frac{196}{15} \zeta_2^2 + \frac{9728}{315} \zeta_3^3, \tag{5.6} \]

and for the gluon form factor

\[ F_3^g |_{N^2} = C_A N_{F}^2 \varepsilon \left( \frac{16823771}{26244} + \frac{9368}{135} \zeta_5 + \frac{5440}{27} \zeta_3 - \frac{30283}{1458} \zeta_2 - \frac{988}{27} \zeta_2 \zeta_3 + \frac{14018}{405} \zeta_2^2 \right) \\
+ C_A^2 N_{F} \varepsilon \left( \frac{196900}{243} - \frac{800}{9} \zeta_5 - \frac{4208}{9} \zeta_3 - \frac{54 \zeta_2}{3} - \frac{112}{3} \zeta_2 \zeta_3 - \frac{2464}{45} \zeta_2^2 \right) \\
+ C_F C_A N_{F} \left( \frac{-10508593}{2916} + \frac{17092}{27} \zeta_5 + \frac{240934}{243} \zeta_3 + \frac{4064}{9} \zeta_3^2 + \frac{8869}{54} \zeta_2^2 \right) \\
+ \frac{640}{9} \zeta_2 \zeta_3 + \frac{28823}{270} \zeta_2^2 + \frac{23624}{315} \zeta_3^2 \right) \\
+ C_F^2 N_{F} \varepsilon \left( \frac{18613}{54} - \frac{3080}{3} \zeta_5 + \frac{10552}{9} \zeta_3 - \frac{272 \zeta_3^2}{3} - \frac{74}{3} \zeta_2^2 \right) \\
- 16 \zeta_2 \zeta_3 + \frac{328}{5} \zeta_2^2 - \frac{35648}{315} \zeta_3^2 \right) \tag{5.7} \]

The UV-renormalization of the form factors is derived in Section 2 above. Applying (2.17) and (2.18) yields the expansion coefficients of the renormalized form factors. These are in the space-like kinematics:

\[ F_3^g = C_F^2 \left[ - \frac{4}{3} e^6 - \frac{6}{e^5} + \frac{1}{e^4} (2 \zeta_2 - 25) - \frac{1}{e^3} \left( 3 \zeta_2 - \frac{100 \zeta_3}{3} + 83 \right) \right. \]
\[ + \frac{1}{e^2} \left( \frac{213}{2} \zeta_2^2 - \frac{77 \zeta_2}{2} + 138 \zeta_3 - \frac{515}{2} \right) \]
\[ + \frac{1}{e} \left( \frac{1461}{10} \zeta_2^3 - \frac{214 \zeta_2 \zeta_3}{3} - \frac{467 \zeta_2}{2} + \frac{2119 \zeta_3}{3} + \frac{644 \zeta_5}{5} - \frac{9073}{12} \right) \]
\[ + \left. \left( - \frac{53675}{24} - \frac{13001 \zeta_2}{12} + \frac{12743 \zeta_3^2}{40} - \frac{9095 \zeta_2^3}{252} + \frac{2669 \zeta_3 + 61 \zeta_3 \zeta_2}{5} \right) \right) \right] \\
+ C_F^2 C_A \left[ - \frac{11}{e^5} - \frac{1}{e^4} \left( \frac{361}{18} + 2 \zeta_2 \right) + \frac{1}{e^3} \left( - \frac{1703}{54} - 26 \zeta_3 + \frac{27 \zeta_2}{2} \right) \right] \]
\[
\begin{align*}
+ & \frac{1}{\epsilon^2} \left( \frac{6820}{81} - \frac{482 \zeta_3}{9} + \frac{1487 \zeta_2}{36} - \frac{83 \zeta_2^2}{5} \right) \\
+ & \frac{1}{\epsilon} \left( \frac{374149}{486} - \frac{142 \zeta_5 + 215 \zeta_3 \zeta_2}{3} - \frac{4151 \zeta_3}{6} + \frac{31891 \zeta_2}{108} - \frac{2975 \zeta_2^2}{72} \right) \\
+ & \left( \frac{11169211}{2916} - \frac{6890 \zeta_5}{9} - \frac{806 \zeta_3 \zeta_2}{3} - \frac{19933 \zeta_3}{6} \\
& \quad + \frac{1616 \zeta_2^2}{3} + \frac{537803 \zeta_2}{324} - \frac{723739 \zeta_2^2}{2160} + \frac{18619 \zeta_2^3}{1260} \right) \\
\end{align*}
\]

\[
\quad + C_F C_A^2 \left[ - \frac{1331}{81 \epsilon^4} + \frac{1}{\epsilon^3} \left( - \frac{2866}{243} - \frac{110 \zeta_3}{27} \right) + \frac{1}{\epsilon^2} \left( \frac{11669}{486} - \frac{902 \zeta_3}{27} + \frac{1625 \zeta_2}{81} - \frac{88 \zeta_2^2}{45} \right) \right. \\
\quad \left. + \frac{1}{\epsilon} \left( - \frac{139345}{8748} - \frac{136 \zeta_5}{3} - \frac{88 \zeta_3 \zeta_2}{9} + \frac{3526 \zeta_3}{27} - \frac{7163 \zeta_2}{243} - \frac{166 \zeta_2^2}{15} \right) \right. \\
\quad \left. + \left( - \frac{51082685}{52488} - \frac{434 \zeta_5}{9} + \frac{416 \zeta_3 \zeta_2}{3} + \frac{505087 \zeta_3}{486} \right) \right] \\
\]

\[
\quad + C_F^2 N_F \left[ \frac{2}{\epsilon^5} + \frac{35}{9 \epsilon^4} + \frac{1}{\epsilon^3} \left( - \frac{139 \zeta_3}{27} + 3 \zeta_2 \right) + \frac{1}{\epsilon^2} \left( - \frac{775}{81} - \frac{110 \zeta_3}{9} - \frac{133 \zeta_2}{18} \right) \right. \\
\quad \left. + \frac{1}{\epsilon} \left( - \frac{24761}{243} + \frac{469 \zeta_3}{27} - \frac{2183 \zeta_2}{54} - \frac{287 \zeta_2^2}{36} \right) \right. \\
\quad \left. + \left( - \frac{691883}{1458} - \frac{386 \zeta_5}{9} + \frac{35 \zeta_3 \zeta_2}{3} + \frac{21179 \zeta_3}{81} - \frac{16745 \zeta_2}{81} - \frac{8503 \zeta_2^2}{1080} \right) \right] \\
\]

\[
\quad + C_F C_A N_F \left[ \frac{484}{81 \epsilon^4} + \frac{1}{\epsilon^3} \left( - \frac{752}{243} + \frac{20 \zeta_2}{27} \right) + \frac{1}{\epsilon^2} \left( - \frac{2068}{243} + \frac{212 \zeta_3}{27} - \frac{476 \zeta_2}{81} \right) \right. \\
\quad \left. + \frac{1}{\epsilon} \left( - \frac{8659}{2187} + \frac{964 \zeta_3}{81} - \frac{2594 \zeta_2}{243} + \frac{44 \zeta_2^2}{15} \right) \right. \\
\quad \left. + \left( \frac{1700171}{6561} - \frac{4 \zeta_5}{3} + \frac{4 \zeta_3 \zeta_2}{3} - \frac{4288 \zeta_3}{27} + \frac{115555 \zeta_2}{729} + \frac{2 \zeta_2^2}{27} \right) \right] \\
\]

\[
\quad + C_F N_F^2 \left[ \frac{44}{81 \epsilon^4} - \frac{8}{243 \epsilon^3} + \frac{1}{\epsilon^2} \left( \frac{46}{81} + \frac{4 \zeta_2}{9} \right) + \frac{1}{\epsilon} \left( \frac{2417}{2187} - \frac{8 \zeta_5}{81} - \frac{20 \zeta_2}{27} \right) \right. \\
\quad \left. + \left( - \frac{190931}{13122} \frac{416 \zeta_3}{243} - \frac{824 \zeta_2}{81} - \frac{188 \zeta_2^2}{135} \right) \right] \\
\]

\[
\quad + C_F N_F V \left[ \left( \frac{N^2 - 4}{N} \right) \left[ 4 - \frac{2 \zeta_2^2}{5} + 10 \zeta_2 + \frac{14 \zeta_5}{3} - \frac{80 \zeta_5}{3} \right] \right], \\
\quad F_3^g = C_A^3 \left[ - \frac{4}{3 \epsilon^6} - \frac{55}{3 \epsilon^3} - \frac{9079}{162 \epsilon^4} + \frac{1}{\epsilon^3} \left( \frac{5453}{486} + \frac{22 \zeta_3}{3} + \frac{77 \zeta_2}{54} \right) \right] \quad (5.8)
\]
The $O(\epsilon)$ contributions to the $N_F$ parts of the UV-renormalized space-like quark and gluon form factors are given by,

\[
F_3^q|_{N_F} = +C_F N_F^2 \epsilon \left( -\frac{3769249}{26244} + \frac{88}{135} \zeta_5 + \frac{2632}{243} \zeta_3 - \frac{5515}{81} \zeta_2 + \frac{8}{3} \zeta_2 \zeta_3 - \frac{952}{81} \zeta_2^2 \right)
\]

\[
+ C_F C_A N_F \epsilon \left( \frac{1552436}{729} - \frac{11596}{45} \zeta_5 - \frac{1214351}{729} \zeta_3 + \frac{3988}{27} \zeta_3^2 + \frac{4933141}{4374} \zeta_2^2 + \frac{1966}{27} \zeta_2 \zeta_3 + \frac{4579}{405} \zeta_2^2 + \frac{2762}{945} \zeta_3^2 \right)
\]
\[ +C_F^2 N_F \epsilon \left( -\frac{15199979}{8748} - \frac{10769}{135} \zeta_5 + \frac{553882}{243} \zeta_3 - \frac{6881}{27} \zeta_2^2 - \frac{961699}{972} \zeta_2 \right. \]
\[ - \frac{4627}{54} \zeta_2 \zeta_3 + \frac{94747}{3240} \zeta_2^2 - \frac{425813}{7560} \zeta_2^3 \left( \frac{N^2 - 4}{N} \right) \epsilon \left( \frac{170}{3} + \frac{752}{9} \zeta_5 + \frac{94}{9} \zeta_3 - \frac{344}{3} \zeta_2^2 + \frac{260}{3} \zeta_2 \right. \]
\[ +30 \zeta_2 \zeta_3 - \frac{196}{15} \zeta_2^2 - \frac{9728}{315} \zeta_2^3 \right), \quad (5.10) \]

and for the gluon form factor

\[ F_3^g |_{N_F} = +C_A N_F \left( \frac{6599393}{26244} + \frac{1844}{135} \zeta_5 + \frac{8396}{81} \zeta_3 - \frac{25315}{1458} \zeta_2 - \frac{172}{27} \zeta_2 \zeta_3 + \frac{2453}{405} \zeta_2^2 \right) \]
\[ +C_A^2 N_F \left( \frac{18825781}{8748} + \frac{1682}{45} \zeta_5 + \frac{270232}{729} \zeta_3 - \frac{6251}{27} \zeta_2^2 + \frac{867919}{4374} \zeta_2 \right. \]
\[ - \frac{881}{9} \zeta_2 \zeta_3 + \frac{33403}{405} \zeta_2^2 + \frac{133627}{7560} \zeta_2^3 \left( \frac{360181}{972} - \frac{224}{9} \zeta_5 - \frac{1960}{9} \zeta_3 - \frac{277}{9} \zeta_2 + \frac{32}{3} \zeta_2 \zeta_3 - \frac{208}{15} \zeta_2^2 \right) \]
\[ +C_F C_A N_F \left( \frac{7017335}{5832} + \frac{7588}{27} \zeta_5 - \frac{92894}{243} \zeta_3 + \frac{4064}{9} \zeta_2^2 + \frac{986}{54} \zeta_2 \right. \]
\[ + \frac{1960}{9} \zeta_2 \zeta_3 - \frac{59987}{540} \zeta_2^2 + \frac{23624}{315} \zeta_2^3 \left( \frac{18613}{54} - \frac{3080}{3} \zeta_5 + \frac{10552}{9} \zeta_3 - \frac{272}{3} \zeta_2^2 - \frac{74}{3} \zeta_2 \right. \]
\[ -16 \zeta_2 \zeta_3 + \frac{328}{5} \zeta_2^2 - \frac{35648}{315} \zeta_2^3 \right). \quad (5.11) \]

6. Infrared pole structure

According to ref. [13, 20], the general infrared pole structure of a renormalised QCD amplitude is related to the ultraviolet behaviour of an effective operator in soft-collinear effective theory. These poles can therefore be subtracted by means of a multiplicative renormalization factor \( Z \). This means that the finite remainders of a scattering amplitude \( M^F \) is obtained from the full amplitude \( M \) via the relation,

\[ M^F = Z^{-1} M. \quad (6.1) \]
In general, the scattering amplitude $M$ and $Z$ are matrices in colour space. However, in the context of the quark and gluon form factors, the colour matrix is trivial. The UV renormalised amplitudes $M$ and $M_F$ have perturbative expansions,

$$ M = 1 + \sum_{i=1} \left( \frac{\alpha_s(\mu^2)}{4\pi} \right)^i M_i, \quad (6.2) $$

$$ M_F = 1 + \sum_{i=1} \left( \frac{\alpha_s(\mu^2)}{4\pi} \right)^i M_i^F, \quad (6.3) $$

while

$$ \log(Z) = \sum_{i=1} \left( \frac{\alpha_s(\mu^2)}{4\pi} \right)^i Z_i. \quad (6.4) $$

We can now solve eq. (6.1) order by order in the strong coupling,

$$ \text{Poles}(M_1) = Z_1, \quad (6.5) $$

$$ \text{Poles}(M_2) = Z_2 + \frac{M_1^2}{2}, \quad (6.6) $$

$$ \text{Poles}(M_3) = Z_3 - \frac{M_1^3}{3} + M_2 M_1, \quad (6.7) $$

$$ \text{Poles}(M_4) = Z_4 + \frac{M_1^4}{4} - M_1^2 M_2 + M_1 M_3 + \frac{M_2^2}{2}, \quad (6.8) $$

$$ \text{Poles}(M_5) = Z_5 - \frac{M_1^5}{5} + M_1^3 M_2 - M_1^2 M_3 - M_1 M_2^2 + M_1 M_4 + M_2 M_3. \quad (6.9) $$

The deepest infrared pole for the $i$-loop amplitude is $\epsilon^{-2i}$. However, the deepest pole in the $Z_i$-factor is $\epsilon^{-i-1}$. All of the deepest poles are obtained directly from the lower loop amplitudes - which must be known to an appropriately high order in $\epsilon$. For example, to obtain the correct pole structure for $M_i$, one needs knowledge of $M_1$ through to $O(\epsilon^{2i-3})$.

We find that the infrared pole structure of the renormalised form factors is given by ($i = q, g$ and $C_q = C_F$, $C_g = C_A$ for the cusp anomalous dimension):

$$ \text{Poles}(F_1^i) = -\frac{C_i \gamma_0^{\text{cusp}}}{2 \epsilon^2} + \frac{\gamma_i^0}{\epsilon}, \quad (6.10) $$

$$ \text{Poles}(F_2^i) = \frac{3C_i \gamma_0^{\text{cusp}} \beta_0}{8 \epsilon^4} + \frac{1}{\epsilon^2} \left( \frac{\beta_0 \gamma_0^i}{2} - \frac{C_i \gamma_1^{\text{cusp}}}{8} \right) + \frac{\gamma_i^0}{2 \epsilon} + \frac{(F_1^i)^2}{2}, \quad (6.11) $$

$$ \text{Poles}(F_3^i) = -\frac{11\beta_0^2 C_i \gamma_0^{\text{cusp}}}{36 \epsilon^4} + \frac{1}{\epsilon^2} \left( \frac{5 \beta_0 C_i \gamma_1^{\text{cusp}}}{36} + \frac{\beta_0^2 \gamma_i^0}{3} + \frac{2 C_i \gamma_2^{\text{cusp}} \beta_1}{9} \right) $$

$$ + \frac{1}{\epsilon^2} \left( \frac{-\beta_0 \gamma_1^i}{3} - \frac{C_i \gamma_2^{\text{cusp}}}{18} - \frac{\beta_1 \gamma_0^i}{3} \right) + \frac{\gamma_i^0}{3 \epsilon} + \frac{(F_1^i)^3}{3} + F_2^q F_1^q. \quad (6.12) $$

Note that the full (all-orders) expressions for $F_2^q$ are recycled on the right-hand-side. The coefficients of the cusp soft anomalous dimension $\gamma_i^{\text{cusp}}$ are known to three-loop order [14].
and are given by:

\[ \gamma_{cusp}^0 = 4, \]
\[ \gamma_{cusp}^1 = C_A \left( \frac{268}{9} - \frac{4\pi^2}{3} \right) - \frac{40N_F}{9}, \]
\[ \gamma_{cusp}^2 = C_A^2 \left( \frac{490}{3} - \frac{536\pi^2}{27} + \frac{44\pi^4}{45} + \frac{88\zeta_3}{3} \right) + C_A N_F \left( -\frac{836}{27} + \frac{80\pi^2}{27} - \frac{112\zeta_3}{3} \right) + C_F N_F \left(-\frac{110}{3} + 32\zeta_3 \right) - 16N_F^2. \]

while the quark and gluon collinear anomalous dimensions \( \gamma^q_i \) and \( \gamma^g_i \) in the conventional dimensional regularisation scheme are also known to three-loop order \([19, 20]\) and are given by:

\[ \gamma^q_0 = -3C_F, \]
\[ \gamma^q_1 = C_F^2 \left(-\frac{3}{2} + 2\pi^2 - 24\zeta_3 \right) + C_F C_A \left(-\frac{961}{54} - \frac{11\pi^2}{6} + 26\zeta_3 \right) + C_F N_F \left(\frac{65}{27} + \frac{\pi^2}{3} \right), \]
\[ \gamma^q_2 = C_F^2 N_F \left(\frac{2953}{54} - \frac{13\pi^2}{9} - \frac{14\pi^4}{27} + \frac{256\zeta_3}{9} \right) + C_F N_F^2 \left(\frac{2417}{729} - \frac{10\pi^2}{27} - \frac{8\zeta_3}{9} \right) + C_F C_A N_F \left(-\frac{8659}{729} + \frac{1297\pi^2}{243} + \frac{11\pi^4}{45} - \frac{964\zeta_3}{27} \right) + C_A C_F^2 \left(-\frac{29}{2} - 3\pi^2 - \frac{8\pi^4}{5} - 68\zeta_3 + \frac{16\pi^2\zeta_3}{3} + 240\zeta_5 \right) + C_A C_F \left(-\frac{151}{4} + \frac{205\pi^2}{9} + \frac{247\pi^4}{135} - \frac{844\zeta_3}{3} + \frac{8\pi^2\zeta_3}{3} - 120\zeta_5 \right) + C_A^2 C_F \left(\frac{139345}{2916} - \frac{7163\pi^2}{486} - \frac{83\pi^4}{90} + \frac{3526\zeta_3}{9} + \frac{44\pi^2\zeta_3}{9} - 136\zeta_5 \right), \]
\[ \gamma^g_0 = -\frac{11C_A}{3} + \frac{2N_F}{3}, \]
\[ \gamma^g_1 = C_A^2 \left(-\frac{692}{27} + \frac{11\pi^2}{18} + 2\zeta_3 \right) + C_A N_F \left(\frac{128}{27} - \frac{\pi^2}{9} \right) + 2C_F N_F, \]
\[ \gamma^g_2 = C_A^3 \left(-\frac{97186}{729} + \frac{6109\pi^2}{486} - \frac{319\pi^4}{270} + \frac{122\zeta_3}{3} - \frac{20\pi^2\zeta_3}{9} - 16\zeta_5 \right) + C_A^2 N_F \left(\frac{30715}{1458} - \frac{599\pi^2}{243} + \frac{41\pi^4}{135} + \frac{356\zeta_3}{27} \right) + C_F C_A N_F \left(\frac{1217}{27} - \frac{\pi^2}{3} - \frac{4\pi^4}{45} - \frac{152\zeta_3}{9} \right) - C_F^2 N_F + C_A N_F^2 \left(-\frac{269}{1458} + \frac{10\pi^2}{81} - \frac{56\zeta_3}{27} \right) - \frac{11C_F N_F^2}{9}. \]
Taking this one step further, we find that the pole structure of the renormalised four-loop quark form factor is given by

\[
Poles(F_i^4) = \frac{25\beta_0^2 C_i \gamma_{0,\text{cusp}}}{96\epsilon^3} + \frac{\beta_0 (24\beta_0^2 \gamma_{0,\text{cusp}} + 13\beta_0 C_i \gamma_{1,\text{cusp}} + 40C_i \gamma_{0,\text{cusp}} \beta_1)}{96\epsilon^4} \\
+ \frac{1}{\epsilon^3} \left( \frac{\beta_0 C_i \gamma_{2,\text{cusp}}}{96} + \frac{3\beta_1 C_i \gamma_{1,\text{cusp}}}{32} + \frac{\beta_0^2 \gamma_{1,\text{cusp}}}{4} + \frac{\beta_1 \beta_0 \gamma_{1,\text{cusp}}}{2} + \frac{5C_i \gamma_{0,\text{cusp}} \beta_2}{32} \right) \\
+ \frac{1}{\epsilon^2} \left( -\frac{\beta_1 \gamma_{1,\text{cusp}}}{4} - \frac{C_i \gamma_{3,\text{cusp}}}{32} - \frac{\beta_0 \gamma_{2,\text{cusp}}}{4} - \frac{\beta_2 \gamma_{1,\text{cusp}}}{4} \right) + \frac{\gamma_{3,\text{cusp}}}{4\epsilon} \\
+ \frac{(F_i^4)^4}{4} + (F_i^4)^2 F_2^i - \frac{(F_3^i)^2}{2} - F_1^i F_3^i. \tag{6.22}
\]

In this expression, we assume Casimir scaling of the cusp anomalous dimension to hold at four loops [20, 21], such that only a universal \( \gamma_{3,\text{cusp}} \) appears. If, contrary to expectations, Casimir scaling should be violated at this order, different \( \gamma_{3,\text{cusp}} \) would appear in the double pole terms of the quark and gluon form factors at four loops.

Eq. (6.22) shows that in order to make use of a calculation of the pole parts of the four-loop form factors to extract the cusp and collinear anomalous dimensions, one requires the finite parts of the three-loop form factor for \( \gamma_{3,\text{cusp}} \), and of the subleading \( \mathcal{O}(\epsilon) \) parts for \( \gamma_{3,\text{cusp}}^q, g \). For all colour-factor contributions proportional to \( N_F \), these are provided in the previous section. The required subleading terms in higher orders in \( \epsilon \) from the one-loop and two-loop form factors were summarized in Section 2 above.

7. Effective Theory Matching Coefficients

It is well known that fixed-order perturbation theory is not reliable for physical quantities involving several disparate scales. In such cases, higher-order corrections are enhanced by large logarithms of scale ratios. Experimentally relevant examples are the Drell-Yan and Higgs production processes in hadron-hadron colliders. When the phase space for soft gluon emission is constrained, large logarithmic threshold corrections appear of the form

\[
\alpha_s^k \left[ \ln^{m-1}(1-z) \right] (1-z), \quad (m \leq 2k),
\]

where \( (1-z) \) is the fraction of centre-of-mass energy of the initial partons available for soft gluon radiation. These spoil the convergence of the perturbative series. The resummation of these so-called Sudakov-logarithms has been accomplished to fourth logarithmic order [7], using the exponentiation properties of the coefficient functions in momentum space [58].

An alternative resummation framework is provided by soft-collinear effective field theory (SCET), which is based on the idea to split the calculation into a series of single-scale problems by successively integrating out the physics associated with the largest remaining scale. The SCET framework [9] originated in the study of heavy quarks, and has been subsequently generalized to massless collider processes [59]. The infrared poles in the high energy theory (QCD) get transformed into ultraviolet poles in the effective theory [8, 13] and can then be resummed by renormalization-group (RG) evolution from the larger scales.
to the smaller ones. Of course the SCET must match precisely onto the high energy theory, and this is achieved by computing matrix elements in both the SCET and QCD and adjusting the Wilson coefficients so that they agree. If the matching is performed on-shell, then the matching coefficients relevant for Drell Yan and Higgs production can be obtained from the quark and gluon form factors respectively. Therefore, we can utilise the results presented in the previous sections to compute the matching conditions through to three-loops. Results up to two loops were obtained previously in [10–12].

The renormalised form-factors are infrared divergent. In the effective field theory, these infrared divergences are transformed into ultraviolet poles. The matching coefficient $C^i (i = q, g)$ is obtained by extracting the poles using a renormalisation factor such that,

$$ C^i(\alpha_s(\mu^2), s_{12}, \mu^2) = \lim_{\epsilon \to 0} Z_i^{-1}(\epsilon, s_{12}, \mu) F^i(\epsilon, s_{12}, \mu^2). $$

(7.2)

The matching coefficients have the perturbative expansion,

$$ C^i(\alpha_s(\mu^2), s_{12}, \mu^2) = 1 + \sum_{n=1}^{\infty} \left( \frac{\alpha_s(\mu^2)}{4\pi} \right)^n C^i_n(s_{12}, \mu^2). $$

(7.3)

They are are known to two loop order for Drell-Yan [10, 11] and Higgs [12] production,

$$ C^q_1 = C_F \left( -L^2 + 3L - 8 + \zeta_2 \right), $$

(7.4)

$$ C^q_2 = C_F^2 \left( \frac{1}{2} L^4 - 3L^3 + \left( \frac{25}{2} - \zeta_2 \right) L^2 + \left( -\frac{45}{2} + 24\zeta_3 - 9\zeta_2 \right)L 
+ \frac{255}{8} - 30\zeta_3 + 21\zeta_2 - \frac{83}{10} \zeta_2^2 \right) 
+ C_F C_A \left( \frac{11}{9} L^3 + \left( -\frac{233}{18} + 2\zeta_2 \right) L^2 + \left( \frac{2545}{54} - 26\zeta_3 + \frac{22}{3} \zeta_2 \right)L 
- \frac{51157}{648} + \frac{313}{9} \zeta_3 - \frac{337}{18} \zeta_2 + \frac{44}{5} \zeta_2^2 \right) 
+ C_F N_F \left( -\frac{2}{9} L^3 + \frac{19}{9} L^2 + \left( -\frac{209}{27} - \frac{4}{3} \zeta_2 \right)L + \frac{4085}{324} + \frac{2}{9} \zeta_3 + \frac{23}{9} \zeta_2 \right), $$

(7.5)

$$ C^g_1 = C_A \left( -L^2 + \zeta_2 \right), $$

(7.6)

$$ C^g_2 = C_A^2 \left( \frac{1}{2} L^4 + \frac{11}{9} L^3 + \left( -\frac{67}{9} + \zeta_2 \right) L^2 + \left( \frac{80}{27} - 2\zeta_3 - \frac{22}{3} \zeta_2 \right)L 
+ \frac{5105}{162} - \frac{143}{9} \zeta_3 + \frac{67}{6} \zeta_2 + \frac{1}{2} \zeta_2^2 \right) 
+ C_A N_F \left( -\frac{2}{9} L^3 + \frac{10}{9} L^2 + \left( \frac{52}{27} + \frac{4}{3} \zeta_2 \right)L - \frac{916}{81} - \frac{46}{9} \zeta_3 - \frac{5}{3} \zeta_2 \right) 
+ C_F N_F \left( 2L - \frac{67}{6} + 8\zeta_3 \right), $$

(7.7)

where $L = \log(-s_{12}/\mu^2)$.
Exploiting the expressions for the renormalised quark and gluon form factors given in eqs. (5.8) and (5.9) respectively, we find that the three-loop matching coefficients are

\[ C_3^q = C_F^3 \left( -\frac{1}{6}L^6 + \frac{3}{2}L^5 + \left( -\frac{17}{2} + \frac{1}{2}\zeta_2 \right)L^4 + \left( 9\zeta_2 + 27 - 24\zeta_3 \right)L^3 \\
+ \left( 102\zeta_3 - \frac{507}{8} - \frac{105}{2}\zeta_2 + \frac{83}{10}\zeta_2^2 \right)L^2 \\
+ \left( -214\zeta_3 - 240\zeta_5 - 8\zeta_2\zeta_3 + \frac{357}{2}\zeta_2 - \frac{207}{10}\zeta_2^2 + \frac{785}{8} \right)L \\
- \frac{413}{5}\zeta_2^2 + 664\zeta_5 - \frac{6451}{24}\zeta_2 + \frac{37729}{630}\zeta_2^2 - 470\zeta_3 + 250\zeta_2\zeta_3 - \frac{2539}{12} + 16\zeta_3^2 \right) \]

\[ + C_F^2 C_A \left( -\frac{11}{9}L^5 + \left( \frac{299}{18} - 2\zeta_2 \right)L^4 + \left( -\frac{2585}{27} + 26\zeta_3 - \frac{1}{9}\zeta_2 \right)L^3 \\
+ \left( \frac{206317}{648} - \frac{1807}{9}\zeta_3 + \frac{502}{9}\zeta_2 - \frac{34}{5}\zeta_2^2 \right)L^2 \\
+ \left( -\frac{13805}{24} + 120\zeta_5 + \frac{2441}{3}\zeta_3 - \frac{11260}{27}\zeta_2 - 10\zeta_2\zeta_3 + \frac{162}{5}\zeta_3^2 \right)L \\
+ \frac{415025}{648} - \frac{2756}{9}\zeta_5 - \frac{18770}{27}\zeta_3 + \frac{296}{3}\zeta_2^2 + \frac{535835}{648}\zeta_2 - \frac{3751}{9}\zeta_2\zeta_3 \\
- \frac{4943}{270}\zeta_2^2 + \frac{12676}{315}\zeta_2^3 \right) \]

\[ + C_F^2 N_F \left( \frac{2}{9}L^5 - \frac{25}{9}L^4 - \left( \frac{410}{27} + \frac{10}{9}\zeta_2 \right)L^3 + \left( -\frac{12815}{324} + \frac{70}{9}\zeta_3 - \frac{112}{9}\zeta_2 \right)L^2 \\
+ \left( \frac{3121}{108} - \frac{610}{9}\zeta_3 + \frac{1618}{27}\zeta_2 + \frac{28}{5}\zeta_2^2 \right)L \\
+ \frac{41077}{972} - \frac{416}{9}\zeta_5 + \frac{13184}{81}\zeta_3 - \frac{31729}{324}\zeta_2 - \frac{38}{9}\zeta_3\zeta_2 - \frac{331}{27}\zeta_2^2 \right) \]

\[ + C_F^2 C_A \left( -\frac{121}{54}L^4 + \left( \frac{2869}{81} - \frac{44}{9}\zeta_2 \right)L^3 + \left( -\frac{18682}{81} + 88\zeta_3 + \frac{26}{9}\zeta_2 - \frac{44}{9}\zeta_2^2 \right)L^2 \\
+ \left( \frac{1045955}{1458} + 136\zeta_5 - \frac{17464}{27}\zeta_3 + \frac{17366}{81}\zeta_2 + \frac{8}{3}\zeta_2\zeta_3 - \frac{94}{3}\zeta_3^2 \right)L \\
- \frac{51082685}{52488} - \frac{434}{9}\zeta_5 + \frac{505087}{486}\zeta_3 - \frac{1136}{9}\zeta_3^2 - \frac{412315}{729}\zeta_2 + \frac{416}{3}\zeta_2\zeta_3 \\
+ \frac{22157}{270}\zeta_2^2 - \frac{6152}{189}\zeta_2^3 \right) \]

\[ + C_F C_A N_F \left( \frac{22}{27}L^4 + \left( \frac{974}{81} + \frac{8}{9}\zeta_2 \right)L^3 + \left( \frac{5876}{81} - 8\zeta_3 + \frac{16}{3}\zeta_2 \right)L^2 \\
+ \left( -\frac{154919}{729} + \frac{724}{9}\zeta_3 - \frac{5864}{81}\zeta_2 + \frac{44}{15}\zeta_2^2 \right)L \\
+ \frac{1700171}{6561} - \frac{4}{3}\zeta_5 - \frac{4288}{27}\zeta_3 + \frac{115555}{729}\zeta_2 + \frac{4}{3}\zeta_2\zeta_3 + \frac{2}{27}\zeta_2^2 \right) \]

\[ + C_F N_F^2 \left( -\frac{2}{27}L^4 + \frac{76}{81}L^3 + \left( -\frac{406}{81} - \frac{8}{9}\zeta_2 \right)L^2 + \left( \frac{9838}{729} + \frac{16}{27}\zeta_3 + \frac{152}{27}\zeta_2 \right)L \right) \]
\[ C_3^g = C_A^3 \left( -\frac{1}{6}L^6 - \frac{11}{9}L^5 + \left( \frac{281}{54} - \frac{3}{2}L \right) L^4 + \left( \frac{11}{3}L^2 - \frac{73}{10}L \right) L^2 \right) \]

and,

\[ C_3^g = C_A^3 \left( \frac{2}{9}L^5 - \frac{8}{27}L^4 + \left( \frac{734}{81} - \frac{2}{3}L \right) L^3 + \left( \frac{377}{27} - \frac{118}{9}L \right) L^2 \right) \]

These matching coefficients allow to perform the three-loop matching of the SCET-based resummation onto the full QCD calculation.

8. Conclusions

In this paper, we described the calculation of the three-loop quark and gluon form factors in detail. Our results confirm earlier expressions obtained by Baikov et al. [41], which we extended by subleading terms in the fermionic corrections.
The form factors are the simplest QCD objects with non-trivial infrared structure. Recent findings on the relation between massless on-shell QCD amplitudes and operators in soft-collinear effective theory [26], combined with constraints from factorization, has led to the conjecture [20] that their pole terms at a given loop level contain all information needed to predict the pole structure of massless on-shell multi-leg amplitudes at the same loop order. In particular, the cusp anomalous dimension can be extracted from the double pole, and the collinear anomalous dimension from the single pole. At a given loop order, finite and subleading terms from lower loop orders are also required. In this respect, the finite terms presented here will be instrumental for the extraction of the four-loop cusp anomalous dimension, while the subleading terms contribute to the four-loop quark and gluon collinear anomalous dimension.

The three-loop form factors are key ingredients for the fourth order (N^{3}\text{LO}) corrections to the inclusive Drell-Yan and Higgs boson production cross sections. The calculation of these, at least in an improvement to the soft approximation [7, 11], could be envisaged in future work. In view of this application, we derived the hard matching coefficients of the SCET operators to this order. Inclusion of these corrections will lead to a further stabilization of the perturbative prediction under scale variations, and are thus important for precision physics at hadron colliders.

Acknowledgements

This research was supported in part by the Swiss National Science Foundation (SNF) under contract 200020-126691, by the Forschungskredit der Universität Zürich, the UK Science and Technology Facilities Council, by the European Commission’s Marie-Curie Research Training Network under contract MRTN-CT-2006-035505 ‘Tools and Precision Calculations for Physics Discoveries at Colliders’, and by the Helmholtz Alliance “Physics at the Terascale”. EWNG gratefully acknowledges the support of the Wolfson Foundation and the Royal Society.

A. Master Integrals for three-loop form factors

In this appendix, we summarize the \(\epsilon\)-expansions of all master integrals needed for the three-loop form factors. Our notation for the integrals follows [37], using a Minkowskian loop integration measure \(d^d k/(2\pi)^d\). All master integrals are defined in Section 4 above.

With the normalization \(S_\Gamma\) defined in (2.19), all \(M\)-loop integrals have an overall factor of

\[
s_{12}^{-n} \left(i S_\Gamma (-s_{12} \pm i 0)^{-t}\right)^M
\]

(A.1)

where \(n\) is fixed by dimensional arguments. Unlike Refs. [37–39], there is no \((-1)^n\) factor. We expand to the required order for the three-loop form factors (which is typically the order where transcendentality 6 first appears).

The one-loop master integral is:

\[
B_{2,1} = \frac{1}{\epsilon} + 2 + 4\epsilon - \epsilon^2 (2\zeta_3 - 8) - \epsilon^3 \left(\frac{6\zeta_3^2}{5} + 4\zeta_3 - 16\right)
\]
\[-\epsilon^4 \left( \frac{12\zeta_2^2}{5} + 8\zeta_3 + 6\zeta_5 - 32 \right)\]
\[-\epsilon^5 \left( \frac{16\zeta_2^3}{7} + \frac{24\zeta_2^2}{5} - 2\zeta_3^2 + 16\zeta_3 + 12\zeta_5 - 64 \right)\]
\[+\epsilon^6 \left( 128 - 18\zeta_7 - 24\zeta_5 - 32\zeta_3 + 4\zeta_3^2 - \frac{48}{5} \zeta_2^2 \zeta_3 - \frac{12}{5} \zeta_2^2 \zeta_3 - \frac{32}{7} \zeta_3^2 \right) + \mathcal{O}(\epsilon^7).\]

(A.2)

The two-loop two-point and three-point master integrals are:

\[B_{3,1} = \frac{1}{16} - \frac{13}{8} - \frac{115\epsilon}{16} + \epsilon^2 \left( \frac{5\zeta_3}{2} - \frac{865}{32} \right) + \epsilon^3 \left( \frac{3\zeta_2^2}{2} + \frac{65\zeta_3}{4} - \frac{5971}{64} \right)\]
\[+\epsilon^4 \left( \frac{39\zeta_2^2}{4} + \frac{575\zeta_3}{8} + \frac{27\zeta_5}{2} - \frac{39193}{128} \right)\]
\[+\epsilon^5 \left( \frac{44\zeta_2^3}{7} + \frac{345\zeta_2^2}{8} + \frac{25\zeta_3^2}{2} + \frac{4325\zeta_3}{16} + \frac{351\zeta_5}{4} - \frac{249355}{256} \right)\]
\[+\epsilon^6 \left( -\frac{1555105}{512} + \frac{165}{2} \zeta_7 + \frac{3105}{8} \zeta_5 + \frac{29855}{32} \zeta_3 - \frac{325}{4} \zeta_3^2 + \frac{2595}{16} \zeta_2^2 - 15\zeta_2^2 \zeta_3 + \frac{286}{7} \zeta_3^2 \right)\]
\[+\mathcal{O}(\epsilon^7),\]

(A.3)

\[B_{4,2} = \frac{1}{\epsilon^2} + \frac{4}{\epsilon} + 12 - \epsilon (4\zeta_3 - 32) - \epsilon^2 \left( \frac{12\zeta_2^2}{5} + 16\zeta_3 - 80 \right)\]
\[-\epsilon^3 \left( \frac{48\zeta_2^2}{5} + 48\zeta_3 + 12\zeta_5 - 192 \right)\]
\[-\epsilon^4 \left( \frac{32\zeta_2^3}{7} + \frac{144\zeta_2^2}{5} - 8\zeta_3^2 + 128\zeta_3 + 48\zeta_5 - 448 \right)\]
\[+\epsilon^5 \left( 1024 - 36\zeta_7 - 144\zeta_5 - 320\zeta_3 + 32\zeta_3^2 - \frac{384}{5} \zeta_2^2 \zeta_3 - \frac{128}{7} \zeta_3^2 \right)\]
\[+\mathcal{O}(\epsilon^6),\]

(A.4)

\[C_{4,1} = \frac{1}{2\epsilon^2} + \frac{5}{2\epsilon} + \left( \zeta_2 + \frac{19}{2} \right) + \epsilon \left( 5\zeta_2 - 4\zeta_3 + \frac{65}{2} \right)\]
\[-\epsilon^2 \left( \frac{6\zeta_2^2}{5} - 19\zeta_2 + 20\zeta_3 - \frac{211}{2} \right)\]
\[-\epsilon^3 \left( \frac{6\zeta_2^2}{5} + 8\zeta_2 \zeta_3 - 65\zeta_2 + 76\zeta_3 + 24\zeta_5 - \frac{665}{2} \right)\]
\[-\epsilon^4 \left( \frac{528\zeta_2^3}{35} + \frac{114\zeta_2^2}{5} + 40\zeta_2 \zeta_3 - 16\zeta_3^2 - 211\zeta_2 + 260\zeta_3 + 120\zeta_5 - \frac{2059}{2} \right)\]
\[+\epsilon^5 \left( \frac{6305}{2} - 156\zeta_7 - 456\zeta_5 - 844\zeta_3 + 80\zeta_3^2 + 665\zeta_2 - 48\zeta_2 \zeta_3 - 152\zeta_2 \zeta_3 - 78\zeta_2^2 \zeta_3 \right)\]
\[+\frac{48}{5} \zeta_2 \zeta_3 - \frac{528}{7} \zeta_3^2 \right) + \mathcal{O}(\epsilon^6),\]

(A.5)

\[C_{6,2} = \frac{1}{\epsilon^4} - \frac{5\zeta_2}{\epsilon^2} - \frac{27\zeta_3}{\epsilon} - 23\zeta_3^2 + \epsilon (48\zeta_2 \zeta_3 - 117\zeta_5) - \epsilon^2 \left( \frac{456\zeta_2^3}{35} - 267\zeta_3^2 \right)\]
\[ + \epsilon^3 \left( 6\zeta_7 + 240\zeta_2\zeta_5 + \frac{1962}{5}\zeta_2^2\zeta_3 \right) + \mathcal{O}(\epsilon^4). \tag{A.6} \]

The \(B_{\ell,i}\)-type and \(C_{\ell,i}\)-type master integrals read at three loops:

\[
B_{4,1} = \frac{1}{36\epsilon} + \frac{71}{216} - \frac{3115\epsilon}{1296} + \epsilon^2 \left( \frac{7\zeta_3}{9} + \frac{109403}{7776} \right) + \epsilon^3 \left( \frac{497\zeta_3}{54} - \frac{7\pi^4}{540} + \frac{3386467}{46656} \right) + \epsilon^4 \left( \frac{21805\zeta_3}{324} - \frac{1295\pi^4}{3240} + \frac{96885467}{279936} \right) + \epsilon^5 \left( \frac{765821\zeta_3}{1944} - \frac{497\zeta_5}{6} - \frac{4361\pi^4}{3888} - \frac{4\pi^6}{243} + \frac{98\zeta_3^2}{9} + \frac{2631913075}{1679616} \right) + \mathcal{O}(\epsilon^6), \tag{A.7} \]

\[
B_{5,1} = -\frac{1}{4\epsilon^2} - \frac{17}{8\epsilon} - \frac{183}{16} + \epsilon \left( 3\zeta_3 - \frac{1597}{32} \right) + \epsilon^2 \left( \frac{51\zeta_3}{2} + \frac{\pi^4}{20} - \frac{12359}{64} \right) + \epsilon^3 \left( \frac{549\zeta_3}{4} + 15\zeta_5 + \frac{17\pi^4}{40} - \frac{88629}{128} \right) + \epsilon^4 \left( \frac{4791\zeta_3}{8} + \frac{255\zeta_5}{2} + \frac{183\pi^4}{80} + \frac{2\pi^6}{63} - \frac{18\zeta_3^2}{256} - \frac{603871}{279936} \right) + \mathcal{O}(\epsilon^5), \tag{A.8} \]

\[
B_{5,2} = -\frac{1}{3\epsilon^2} - \frac{10}{3\epsilon} - \frac{64}{3} + \epsilon \left( -112 + \frac{22\zeta_3}{3} \right) + \epsilon^2 \left( -528 + \frac{220\zeta_3}{3} + \frac{11\pi^4}{90} \right) + \epsilon^3 \left( -2336 + \frac{1408\zeta_3}{3} + 70\zeta_5 + \frac{11\pi^4}{9} \right) + \epsilon^4 \left( \frac{352\pi^4}{45} + \frac{2464\zeta_3}{3} + 700\zeta_5 - \frac{29824}{3} + \frac{94\pi^6}{567} - \frac{242\zeta_3^2}{3} \right) + \mathcal{O}(\epsilon^5), \tag{A.9} \]

\[
B_{6,1} = \frac{1}{\epsilon^3} + \frac{6}{\epsilon^2} + \frac{24}{\epsilon} + \left( 80 - 6\zeta_3 \right) + \epsilon \left( 240 - 36\zeta_3 - \frac{\pi^4}{10} \right) + \epsilon^2 \left( 672 - 144\zeta_3 - 18\zeta_5 - \frac{3\pi^4}{5} \right) + \epsilon^3 \left( 1792 - 480\zeta_3 - 108\zeta_5 - \frac{12\pi^4}{5} - \frac{2\pi^6}{63} + 18\zeta_3^2 \right) + \mathcal{O}(\epsilon^4), \tag{A.10} \]

\[
B_{6,2} = \frac{1}{3\epsilon^3} + \frac{7}{3\epsilon^2} + \frac{31}{3\epsilon} + \left( \frac{8\zeta_3}{3} + \frac{103}{3} \right) + \epsilon \left( \frac{235}{3} + \frac{56\zeta_3}{3} + \frac{2\pi^4}{45} \right) + \epsilon^2 \left( \frac{19}{3} + 120\zeta_5 + \frac{320\zeta_3}{3} + \frac{14\pi^4}{45} \right) + \epsilon^3 \left( \frac{-3953}{3} + 840\zeta_5 + \frac{1832\zeta_3}{3} + \frac{16\pi^4}{9} + \frac{176\pi^6}{567} - \frac{292\zeta_3^2}{3} \right) + \mathcal{O}(\epsilon^4), \tag{A.11} \]

\[
B_{8,1} = 20\zeta_5 + \epsilon \left( 68\zeta_3^2 + 40\zeta_5 + \frac{10\pi^6}{189} \right) + \epsilon^2 \left( 136\zeta_5^2 + \frac{34\pi^4\zeta_3}{15} + 80\zeta_5 + \frac{20\pi^6}{189} + 450\zeta_7 \right) + \mathcal{O}(\epsilon^3), \tag{A.12} \]
The genuine three-loop vertex integrals are:

\[ C_{6,1} = \frac{1}{2\epsilon^3} + \frac{7}{2\epsilon^2} + \frac{1}{\epsilon} \left( \frac{\pi^2}{6} + \frac{33}{2} \right) + \left( \frac{7\pi^2}{6} - \frac{5\zeta_3}{3} + \frac{131}{2} \right) + \epsilon \left( \frac{11\pi^2}{2} - \frac{35\zeta_3}{3} - \frac{\pi^4}{20} + \frac{473}{2} \right) + \epsilon^2 \left( \frac{131\pi^2}{6} - \frac{5\pi^2\zeta_3}{3} - \frac{165\zeta_3}{3} - \frac{7\pi^4}{20} + \frac{1611}{2} \right) + \epsilon^3 \left( \frac{473\pi^2}{6} - \frac{35\pi^2\zeta_3}{3} - \frac{655\zeta_3}{3} - \frac{33\pi^4}{20} - \frac{61\pi^6}{756} + \frac{25\zeta_3^2}{2} + \frac{5281}{2} \right) + \mathcal{O}(\epsilon^4), \tag{A.13} \]

\[ C_{8,1} = \frac{1}{\epsilon^3} + \frac{2}{\epsilon^4} + \frac{1}{\epsilon^3} \left( \frac{5\pi^2}{6} + 4 \right) + \frac{1}{\epsilon^2} \left( 8 - \frac{5\pi^2}{3} - 29\zeta_3 \right) + \frac{1}{\epsilon} \left( 16 - \frac{10\pi^2}{3} - 58\zeta_3 - \frac{121\pi^4}{180} \right) + \left( 32 - \frac{20\pi^2}{3} + \frac{29\pi^2\zeta_3}{3} - \frac{116\zeta_3}{3} - \frac{121\pi^4}{90} \right) + \epsilon \left( \frac{58\pi^2\zeta_3}{3} - 232\zeta_3 - 246\zeta_5 - \frac{40\pi^2}{3} + \frac{323\zeta_3^2}{2} - \frac{121\pi^4}{45} + 64 + \frac{163\pi^6}{3780} \right) + \mathcal{O}(\epsilon^2). \tag{A.14} \]

The genuine three-loop vertex integrals are:

\[ A_{5,1} = \frac{1}{24\epsilon^2} + \frac{19}{48\epsilon} + \left( \frac{233}{96} + \frac{\pi^2}{24} \right) + \epsilon \left( \frac{2363}{192} + \frac{19\pi^2}{48} - \frac{11\zeta_3}{12} \right) + \epsilon^2 \left( \frac{7227}{128} + \frac{233\pi^2}{96} + \frac{\pi^4}{80} - \frac{209\zeta_3}{24} \right) + \epsilon^3 \left( \frac{62641}{256} + \frac{2363\pi^2}{192} + \frac{19\pi^4}{160} - \frac{2563\zeta_3}{48} - \frac{11\pi^2\zeta_3}{12} - \frac{35\zeta_5}{4} \right) + \epsilon^4 \left( \frac{1575481}{1536} + \frac{7227\pi^2}{128} + \frac{233\pi^4}{320} - \frac{919\pi^6}{45360} - \frac{25993\zeta_3}{96} - \frac{209\pi^2\zeta_3}{24} + \frac{121\zeta_3^2}{12} - \frac{665\zeta_5}{8} \right) + \mathcal{O}(\epsilon^5), \tag{A.15} \]

\[ A_{5,2} = -\frac{1}{6\epsilon^2} - \frac{5}{3\epsilon} + \left( -\frac{32}{3} - \frac{\pi^2}{12} \right) + \epsilon \left( -56 - \frac{5\pi^2}{6} + \frac{11\zeta_3}{3} \right) + \epsilon^2 \left( -264 - \frac{16\pi^2}{3} + \frac{19\pi^4}{720} + \frac{110\zeta_3}{3} \right) + \epsilon^3 \left( -1168 - 28\pi^2 + \frac{19\pi^4}{72} + \frac{704\zeta_3}{3} + \frac{11\pi^2\zeta_3}{6} + 35\zeta_5 \right) + \epsilon^4 \left( -\frac{14912}{3} - 132\pi^2 + \frac{76\pi^4}{45} + \frac{9011\pi^6}{90720} + 1232\zeta_3 + \frac{55\pi^2\zeta_3}{3} + \frac{121\zeta_3^2}{3} + 350\zeta_5 \right) + \mathcal{O}(\epsilon^5), \tag{A.16} \]

\[ A_{6,1} = \frac{1}{3\epsilon^3} + \frac{8}{3\epsilon^2} + \frac{1}{\epsilon} \left( \frac{44}{3} + \frac{\pi^2}{3} \right) + \left( \frac{8\pi^2}{3} + \frac{208}{3} - \frac{16\zeta_3}{3} \right) \]
\[ A_{6,2} = \frac{2\zeta_3}{\epsilon} + \left( \frac{7\pi^4}{180} + 18\zeta_3 \right) \]
\[ + \epsilon \left( \frac{7\pi^4}{20} + 122\zeta_3 - \frac{2\pi^2\zeta_3}{3} + 10\zeta_5 \right) \]
\[ + \epsilon^2 \left( \frac{427\pi^4}{180} - \frac{163\pi^6}{7560} + 738\zeta_3 - 6\pi^2\zeta_3 - 76\zeta_3^2 + 90\zeta_5 \right) + \mathcal{O}(\epsilon^3), \quad (A.17) \]

\[ A_{6,3} = \frac{1}{6\epsilon^5} + \frac{3}{2\epsilon^2} + \frac{1}{\epsilon^3} \left( 1 - \frac{\pi^2}{6} \right) + \frac{1}{\epsilon^4} \left( 2 - \frac{\pi^2}{3} - 10\zeta_3 \right) \]
\[ + \epsilon \left( \frac{1351}{6} \zeta_3^2 + \frac{55\pi^2}{18} - \frac{\pi^4}{6} + \frac{95\zeta_3}{2} + 3\zeta_5 \right) \]
\[ + \epsilon^2 \left( \frac{2023}{18} + \frac{95\pi^2}{10} + \frac{935\zeta_3}{3} - \frac{3\pi^2\zeta_3}{3} - \frac{10\pi^2\zeta_3}{3} - 36\zeta_5 \right) \]
\[ + \epsilon^3 \left( \frac{26335}{2} + \frac{1351\pi^2}{18} + \frac{11\pi^4}{6} - \frac{7\pi^6}{54} - 1615\zeta_3 - 30\zeta_3^2 + \frac{68\zeta_3^2}{3} - 585\zeta_5 \right) \]
\[ + \mathcal{O}(\epsilon^4), \quad (A.18) \]

\[ A_{7,1} = \frac{1}{4\epsilon^5} + \frac{1}{2\epsilon^4} + \frac{1}{\epsilon^5} \left( 1 - \frac{\pi^2}{6} \right) + \frac{1}{\epsilon^6} \left( 2 - \frac{\pi^2}{3} - 10\zeta_3 \right) \]
\[ + \epsilon \left( 4 - \frac{2\pi^2}{3} - \frac{11\pi^4}{45} - 20\zeta_3 \right) \]
\[ + \epsilon^2 \left( -\frac{22\pi^4}{45} + \frac{4\pi^2}{3} + \frac{14\pi^2\zeta_3}{3} + 8 - 40\zeta_3 - 88\zeta_5 \right) \]
\[ + \epsilon^3 \left( 16 - \frac{8\pi^2}{3} - \frac{44\pi^4}{45} - \frac{943\pi^6}{7560} - 80\zeta_3 + \frac{28\pi^2\zeta_3}{3} + 196\zeta_3^2 - 176\zeta_5 \right) \]
\[ + \mathcal{O}(\epsilon^4), \quad (A.19) \]

\[ A_{7,2} = \frac{\pi^2}{12\epsilon^3} + \frac{1}{\epsilon^2} \left( \frac{\pi^2}{6} + 2\zeta_3 \right) + \frac{1}{\epsilon^3} \left( \frac{\pi^2}{3} + \frac{83\pi^4}{720} + 4\zeta_3 \right) \]
\[ + \epsilon \left( \frac{2\pi^2}{3} + \frac{83\pi^4}{360} + 8\zeta_3 - \frac{5\pi^2\zeta_3}{3} + 15\zeta_5 \right) \]
\[ + \epsilon^2 \left( \frac{4\pi^2}{180} + \frac{83\pi^4}{90720} + 16\zeta_3 - \frac{10\pi^2\zeta_3}{3} - 73\zeta_3^2 + 30\zeta_5 \right) \]
\[ + \mathcal{O}(\epsilon^2), \quad (A.20) \]

\[ A_{7,3} = \frac{1}{\epsilon} \left( -\frac{\pi^2\zeta_3}{6} - 10\zeta_5 \right) + \left( \frac{119\pi^6}{2160} - \frac{31\zeta_3^2}{2} \right) + \mathcal{O}(\epsilon), \quad (A.21) \]
The most complicated three-loop vertex integrals are the nine-line master integrals [39, 40]:

\[ A_{7,4} = \frac{6\zeta_3}{\epsilon^2} + \frac{1}{\epsilon} \left( \frac{11\pi^4}{90} + 36\zeta_3 \right) + \left( 216\zeta_3 - 2\pi^2\zeta_3 + \frac{11\pi^4}{15} + 46\zeta_5 \right) + \epsilon \left( \frac{22\pi^4}{5} - \frac{19\pi^6}{270} + 1296\zeta_3 - 12\pi^2\zeta_3 - 282\zeta_3^2 + 276\zeta_5 \right) + \mathcal{O}(\epsilon^2) , \]  

(A.23)

\[ A_{7,5} = + \left( 2\pi^2\zeta_3 + 10\zeta_5 \right) + \epsilon \left( \frac{11\pi^6}{162} + 12\pi^2\zeta_3 + 18\zeta_3^2 + 60\zeta_5 \right) + \mathcal{O}(\epsilon^2) , \]  

(A.24)

\[ A_{8,1} = \frac{8\zeta_3}{3\epsilon^2} + \frac{1}{\epsilon} \left( \frac{5\pi^4}{27} + 8\zeta_3 \right) + \left( \frac{5\pi^4}{9} - 24\zeta_3 + \frac{52\pi^2\zeta_3}{9} - \frac{352\zeta_5}{3} \right) + \epsilon \left( \frac{-5\pi^4}{3} - \frac{1709\pi^6}{8505} + 72\zeta_3 - \frac{52\pi^2\zeta_3}{3} + \frac{332\zeta_3^2}{3} + 352\zeta_5 \right) + \mathcal{O}(\epsilon^2) . \]  

(A.25)

The most complicated three-loop vertex integrals are the nine-line master integrals [39, 40]:

\[ A_{9,1} = \frac{1}{18\epsilon^5} + \frac{1}{2\epsilon^4} + \frac{1}{\epsilon^3} \left( \frac{-53}{18} - \frac{4\pi^2}{27} \right) + \epsilon \left( \frac{29}{2} + \frac{22\pi^2}{27} - 2\zeta_3 \right) + \epsilon \left( \frac{-8\pi^2}{3} + \frac{158\zeta_3}{9} - \frac{20\pi^4}{81} - \frac{129}{2} \right) + \epsilon \left( \frac{322\pi^4}{405} + 6\pi^2 - \frac{14\pi^2\zeta_3}{3} + \frac{537}{2} - \frac{578\zeta_3}{9} - \frac{238\zeta_5}{3} \right) + \mathcal{O}(\epsilon^2) , \]  

(A.26)

\[ A_{9,2} = \frac{2}{9\epsilon^6} + \frac{5}{6\epsilon^5} + \frac{1}{\epsilon^4} \left( \frac{-20}{9} - \frac{7\pi^2}{27} \right) + \frac{1}{\epsilon^3} \left( \frac{50}{9} - \frac{17\pi^2}{27} - \frac{91\zeta_3}{9} \right) + \frac{1}{\epsilon^2} \left( \frac{4\pi^2}{3} - \frac{166\zeta_3}{9} - \frac{373\pi^4}{1080} - \frac{110}{9} \right) + \epsilon \left( \frac{494\zeta_3}{9} + \frac{179\pi^2\zeta_3}{27} - \frac{16\pi^2\zeta_3}{9} - \frac{16\pi^4}{540} + \frac{170}{9} \right) + \frac{1}{\epsilon} \left( \frac{130}{9} - \frac{32\pi^2}{9} - \frac{1466\zeta_3}{9} + \frac{679\pi^4}{540} + \frac{682\pi^2\zeta_3}{27} - \frac{1300\zeta_5}{3} - \frac{5979\pi^6}{136080} + \frac{169\zeta_3^2}{9} \right) + \mathcal{O}(\epsilon) , \]  

(A.27)

\[ A_{9,4} = \frac{1}{9\epsilon^6} + \frac{8}{9\epsilon^5} + \frac{1}{\epsilon^4} \left( -1 - \frac{10\pi^2}{27} \right) + \frac{1}{\epsilon^3} \left( -\frac{14}{9} - \frac{47\pi^2}{27} - 12\zeta_3 \right) + \frac{1}{\epsilon^2} \left( 17 + \frac{71\pi^2}{27} - \frac{200\zeta_3}{3} - \frac{47\pi^4}{810} \right) + \epsilon \left( -84 - \pi^2 + \frac{940\zeta_3}{9} - \frac{671\pi^4}{540} + \frac{652\pi^2\zeta_3}{27} - \frac{692\zeta_5}{9} \right) \]
\begin{equation} 
\left( 339 - 15\pi^2 - \frac{448\zeta_3}{9} + \frac{2689\pi^4}{1620} + \frac{2188\pi^2\zeta_3}{27} - \frac{524\zeta_5}{50260} + \frac{53563\pi^6}{102060} + \frac{4564\zeta_3^2}{9} \right)
+ \mathcal{O}(\epsilon). \tag{A.28} 
\end{equation}

For the corresponding integrals with homogeneous transcendentality, one finds:

\begin{align*}
A_{9,1n} &= \frac{1}{36\epsilon^6} + \frac{\pi^2}{18\epsilon^4} + \frac{14\zeta_3}{9\epsilon^3} + \frac{47\pi^4}{405\epsilon^2} \\
&\quad + \left( \frac{85}{27}\pi^2\zeta_3 + 20\zeta_5 \right) \frac{1}{\epsilon} + \frac{1160\pi^6}{5103} + \frac{137}{3}\zeta_3^2 + \mathcal{O}(\epsilon) \tag{A.29} \\
A_{9,2n} &= \frac{2}{9\epsilon^6} - \frac{7\pi^2}{27\epsilon^4} - \frac{91\zeta_3}{9\epsilon^3} - \frac{373\pi^4}{1080\epsilon^2} \\
&\quad + \left( \frac{179}{27}\pi^2\zeta_3 - 167\zeta_5 \right) \frac{1}{\epsilon} - \frac{59797}{136080}\pi^6 + \frac{169}{9}\zeta_3^2 + \mathcal{O}(\epsilon) \tag{A.30} \\
A_{9,4n} &= -\frac{1}{9\epsilon^6} + \frac{10\pi^2}{27\epsilon^4} + \frac{12\zeta_3}{\epsilon^3} + \frac{47\pi^4}{810\epsilon^2} \\
&\quad + \left( -\frac{652}{27}\pi^2\zeta_3 + \frac{692}{9}\zeta_5 \right) \frac{1}{\epsilon} - \frac{53563}{102060}\pi^6 - \frac{4564}{9}\zeta_3^2 + \mathcal{O}(\epsilon) \tag{A.31}
\end{align*}

B. Form factors in terms of master integrals

The unrenormalised three-loop form factors can be expressed as a linear combination of master integrals. In the colour factor decomposition as defined in (5.1) and (5.2), these coefficients read:

\begin{equation} 
X_{CF}^q = 
\begin{pmatrix}
489406D^3 & 43304589D^2 & 615952127D & 34015 & 109222498 \\
625 & 3125 & 7500 & 4(2D - 7) & 75(2D - 9)
\end{pmatrix}
\begin{pmatrix}
-50720 & -6816654 & 89728 & 6489724 \\
9(3D - 10) & 11(3D - 14) & 12(2D - 3) & 15D - 4
\end{pmatrix}
\begin{pmatrix}
19326056092 & 7734375(5D - 16) & 13564 & 3(2D - 7)^2 - 12(2D - 3)^2 \\
725 & 1187096 & 745376 & 91648 & 53258146831 & 12(D - 3)^2
\end{pmatrix}
\begin{pmatrix}
54568D^3 & 16060301D^2 & 135964099D & 34015 & 59657807 \\
625 & 9375 & 11250 & 2(2D - 7) & 1300(2D - 9)
\end{pmatrix}
\begin{pmatrix}
36208 & 142784 & 106 & 770008 & 3481535536 \\
27(3D - 10) & 75(2D - 2) & 3(2D - 3) & 15(2D - 4) & 3046875(5D - 16)
\end{pmatrix}
\begin{pmatrix}
-78125(5D - 18) & 234375(5D - 22) & 83104 & 12800 & 26112 \\
35332079719 & 1687500 & 32265012416 & 12800 & 26112
\end{pmatrix}
\begin{pmatrix}
87316D^3 & 8532244D^2 & 15436454D & 418 & 3273151 \\
1875 & 9375 & 3125 & (2D - 7) & 325(2D - 9)
\end{pmatrix}
\end{equation}
\[-C_{6,1} \left( \frac{7D^3}{8} - \frac{1109D^2}{48} + \frac{29395D}{288} + \frac{8475}{64(2D - 7)} + \frac{200}{27(3D - 8)} \right) \]
\[-A_{7,1} \left( \frac{21D^3}{50} - \frac{5907D^2}{500} + \frac{523857D}{5000} - \frac{1213}{12(2D - 7)} + \frac{624(2D - 9)}{10000} \right) \]
\[+ \frac{3}{3(D - 2)} \left( \frac{15D^3}{16} - \frac{733D^2}{32} + \frac{228267D}{1600} + \frac{42745}{2(2D - 7)} + \frac{232399}{8320(2D - 9)} \right) \]
\[+ \frac{1}{45(D - 2)} \left( \frac{601D^3}{1250} - \frac{60199D^2}{6250} + \frac{760189D}{6250} - \frac{1213}{12(2D - 7)} + \frac{2340(2D - 9)}{105625} \right) \]
\[+ \frac{1}{(D - 2)} \left( \frac{21}{15625} \right) \]
\[-A_{7,4} \left( \frac{2489D^3}{5000} - \frac{686707D^2}{5000} + \frac{7042751D}{60000} + \frac{865}{72(2D - 7)} + \frac{235409}{2880(2D - 9)} \right) \]
\[+ \frac{1}{45(D - 2)} \left( \frac{556}{420(D - 3)} - \frac{4397}{5(D - 4)} + \frac{16}{93131696} + \frac{277480707}{100000} \right) \]
\[+ \frac{1}{4921875(5D - 22)} \left( \frac{1}{3(D - 3)^2} \right) \]
\[-A_{7,5} \left( \frac{2489D^3}{5000} - \frac{686707D^2}{5000} + \frac{7042751D}{60000} + \frac{865}{72(2D - 7)} + \frac{235409}{2880(2D - 9)} \right) \]
\[+ \frac{1}{45(D - 2)} \left( \frac{556}{420(D - 3)} - \frac{4397}{5(D - 4)} + \frac{16}{93131696} + \frac{277480707}{100000} \right) \]
\[+ \frac{1}{4921875(5D - 22)} \left( \frac{1}{3(D - 3)^2} \right) \]
\[+ A_{8,1} \left( \frac{341D^3}{80000} - \frac{758793D^2}{80000} + \frac{3243781D}{32000} + \frac{33573}{1024(2D - 7)} + \frac{32}{14(2D - 7)} \right) \]
\[+ \frac{1}{448(D - 3)} \left( \frac{4}{5(D - 4)} \right) \]
\[-B_{8,1} \left( D^3 - 20D^2 + 104D - 176)(D^2 - 7D + 16) \right) \]
\[-C_{8,1} \left( D^3 - 20D^2 + 104D - 176)(D^2 - 7D + 16) \right) \]
\[+ A_{9,1} \left( \frac{243D^3}{1250} - \frac{14661D^2}{3125} + \frac{257769D}{6250} + \frac{256}{5(D - 2)} - \frac{225}{14(D - 3)} \right) \]
\[
- \frac{48}{5(D - 4)} + \frac{4083992}{78125(5D - 16)} + \frac{6463488}{78125(5D - 18)} + \frac{2127008}{546875(5D - 22)} - \frac{4086513}{31250} \\
- A_{9,2} \\
- \frac{3(3D - 14)(D^6 - 41D^5 + 661D^4 - 4992D^3 + 19276D^2 - 37104D + 28288)}{10(5D - 16)(5D - 18)(5D - 22)(D - 3)} \\
- A_{9,4} \\
- \frac{567D^3}{8000} - \frac{125091D^2}{80000} + \frac{1808937D}{1600000} + \frac{2304(2D - 7)}{998400(2D - 9)} + \frac{4067}{574016} - \frac{232399}{7557808} \\
- \frac{16}{75(D - 2)} - \frac{448(D - 3)}{3046875(5D - 16)} - \frac{574016}{234375(5D - 18)} - \frac{7557808}{4921875(5D - 22)} \\
- \frac{38866491}{16000000}
\]

\[
X^q_{C_3^2C_4} = \\
+ B_{4,1} \left( + \frac{225717D^3}{250} - \frac{995657D^2}{625} + \frac{893831341D}{9375} + \frac{45685}{24(2D - 7)} - \frac{990312631}{975(2D - 9)} \right) \\
+ \frac{116080}{21(D - 10)} - \frac{707967}{(3D - 14)} - \frac{371482}{623(5D - 7)} - \frac{125032}{25(D - 2)} - \frac{875}{4(D - 3)} \\
+ \frac{172908907}{405(D - 4)} - \frac{5622111978}{2234375(5D - 16)} - \frac{345702357}{3125(5D - 18)} - \frac{5032168544}{46875(5D - 22)} - \frac{1280}{3(D - 2)^2} \\
+ \frac{2295056}{9} - \frac{3875224}{4(D - 3)} + \frac{2622656}{45(D - 4)^2} - \frac{340096}{45(D - 4)^3} + \frac{15(D - 4)^5}{15(D - 4)^5} \\
- \frac{2728978211}{25000} \\
+ A_{5,1} \left( + \frac{71304D^3}{625} - \frac{12785517D^2}{625} + \frac{156484099D}{12500} - \frac{17815}{3(2D - 7)} - \frac{97578701}{2600(2D - 9)} \right) \\
+ \frac{73232}{63(3D - 10)} + \frac{2467480}{27027(5D - 1)} + \frac{108736}{75(5D - 2)} - \frac{1}{(D - 3)} + \frac{135(D - 4)}{10782693376} \\
+ \frac{9290021216}{11171875(5D - 16)} + \frac{26729196704}{1015625(5D - 18)} - \frac{50176}{78125(5D - 22)} - \frac{179840}{9(D - 4)^2} - \frac{130816}{15(D - 4)^3} \\
+ \frac{23118311953}{15(D - 4)^4} - \frac{1125000}{50176} \\
- A_{5,2} \left( + \frac{25782D^3}{625} - \frac{24183202D^2}{28125} + \frac{437583827D}{84375} + \frac{2036}{3(2D - 7)} + \frac{7969573}{650(2D - 9)} \right) \\
+ \frac{190048}{567(3D - 10)} - \frac{100580}{3(3D - 14)} + \frac{16}{(3D - 8)} - \frac{81081(D - 1)}{1088502038} + \frac{17815}{75(D - 2)} \\
- \frac{1092208}{405(D - 4)} - \frac{9583273136}{33515625(5D - 16)} - \frac{4168502038}{1015625(5D - 18)} - \frac{234375(5D - 22)}{3(D - 2)^2} \\
- \frac{523712}{135(D - 4)^2} - \frac{90176}{45(D - 4)^3} - \frac{35072}{45(D - 4)^4} - \frac{21872512759}{2531250} \\
+ B_{5,1} \left( + \frac{114279D^3}{10000} - \frac{22958287D^2}{10000} + \frac{448177891D}{20000} + \frac{177975}{128(2D - 7)} - \frac{326249}{26(2D - 9)} \right) \\
+ \frac{6128}{15(D - 2)} - \frac{2}{(D - 3)} + \frac{35294}{5(D - 4)} - \frac{25589872}{3046875(5D - 16)} + \frac{831628448}{78125(5D - 18)}
\[
\begin{align*}
- B_{5, 2} &= \left(\frac{+44339D^3}{625} - \frac{4405736D^2}{3125} + \frac{287753909D}{28125} - \frac{237272}{13(2D - 9)} + \frac{16}{27(3D - 10)} \right) + 16 \\
+ A_{6, 1} &= \left(\frac{4678D^3}{625} - \frac{609667D^2}{3125} + \frac{2588633D}{18750} + \frac{10085}{9(3D - 14)} + \frac{54D - 1}{54(D - 1)} \right) + 545 \\
- A_{6, 2} &= \left(\frac{967D^3}{1200} - \frac{1150061D^2}{6000} + \frac{27382277D}{18000} - \frac{509}{8(2D - 7)} + \frac{15600(2D - 9)}{20493136} \right) - 32 \\
+ A_{6, 3} &= \left(\frac{+60679D^3}{2500} - \frac{3322501D^2}{75000} + \frac{436696447D}{150000} + \frac{2545}{12(2D - 7)} + \frac{14641137}{10400(2D - 9)} \right) - 640 \\
- B_{6, 2} &= \left(\frac{+D^3}{16} + \frac{71D^2}{32} - \frac{283D}{64} - \frac{8475}{128(2D - 7)} - \frac{10678}{3D - 8} \right) + 8 \\
- C_{6, 1} &= \left(\frac{+D^3}{16} + \frac{71D^2}{32} - \frac{283D}{64} - \frac{8475}{128(2D - 7)} - \frac{10678}{3D - 8} \right) + 8 \\
+ A_{7, 1} &= \left(\frac{27D^3}{50} - \frac{7000D^2}{500} + \frac{584559D}{5000} - \frac{301}{6(2D - 7)} + \frac{21185}{624(2D - 9)} \right) + 21185 \\
+ A_{7, 2} &= \left(\frac{517D^3}{800} - \frac{142459D^2}{8000} + \frac{9023129D}{80000} + \frac{42745}{192(2D - 7)} + \frac{697197}{16640(2D - 9)} \right) + 21185
\end{align*}
\]
\[
\begin{align*}
&\frac{-112}{143(D - 1)} - \frac{76}{15(D - 2)} - \frac{112}{(D - 4)} - \frac{106444064}{446875(5D - 16)} + \frac{9644612}{40625(5D - 18)} - \frac{33226167}{160000} \\
&- A_{7,3} \left( + \frac{601D^3}{1250} - \frac{29899D^2}{544} + \frac{152417D}{3125} - \frac{301}{6(2D - 7)} + \frac{4237}{468(2D - 9)} \right) \\
&\quad + \frac{54}{3(2D - 2)} + \frac{3683}{210(D - 3)} + \frac{10424832}{1015625(5D - 16)} + \frac{54045184}{703125(5D - 18)} - \frac{12403904}{546875(5D - 22)} \\
&\quad - \frac{3}{(D - 3)^2} - \frac{62500}{15077947} \\
&+ A_{7,4} \left( + \frac{19D^3}{80} - \frac{458683D^2}{60000} + \frac{40603349D}{600000} - \frac{32(2D - 7)}{24} + \frac{5481191}{46026288} \right) \\
&\quad + \frac{118}{15(D - 2)} - \frac{26393}{840(D - 3)} - \frac{325}{5(D - 4)} + \frac{5481191}{203125(5D - 16)} - \frac{140625(5D - 18)}{140625(5D - 18)} \\
&\quad - \frac{328125(5D - 22)}{328125(5D - 22)} - \frac{62067409}{40000} \\
&+ A_{7,5} \left( + \frac{19D^3}{80} - \frac{458683D^2}{60000} + \frac{40603349D}{600000} - \frac{32(2D - 7)}{24} + \frac{5481191}{46026288} \right) \\
&\quad + \frac{118}{15(D - 2)} - \frac{26393}{840(D - 3)} - \frac{325}{5(D - 4)} + \frac{5481191}{203125(5D - 16)} - \frac{140625(5D - 18)}{140625(5D - 18)} \\
&\quad - \frac{328125(5D - 22)}{328125(5D - 22)} - \frac{62067409}{40000} \\
&- A_{8,1} \left( + \frac{1197D^3}{160000} - \frac{979611D^2}{1600000} + \frac{8338443D}{640000} + \frac{29027}{2048(2D - 7)} + \frac{160}{3(D - 2)} \right) \\
&\quad - \frac{13665}{896(D - 3)} - \frac{34}{5(D - 4)} + \frac{23195948}{234375(5D - 16)} - \frac{3147936}{78125(5D - 18)} - \frac{1018736}{546875(5D - 22)} \\
&\quad + \frac{3563}{128(2D - 7)^2} - \frac{206584391}{32000000} \\
&+ B_{8,1} \left( \frac{(D^3 - 20D^2 + 104D - 176)(D^2 - 7D + 16)}{16(2D - 7)(D - 4)} \right) \\
&+ C_{8,1} \left( \frac{(D^3 - 20D^2 + 104D - 176)(D^2 - 7D + 16)}{16(2D - 7)(D - 4)} \right) \\
&- A_{9,1} \left( + \frac{243D^3}{1250} - \frac{14661D^2}{3125} + \frac{257769D}{6250} + \frac{256}{5(D - 2)} - \frac{225}{14(D - 3)} \right) \\
&\quad - \frac{48}{5(D - 4)} + \frac{4083992}{78125(5D - 16)} + \frac{646348}{78125(5D - 18)} + \frac{2127008}{546875(5D - 22)} - \frac{4086513}{31250} \\
&+ A_{9,2} \left( \frac{3(D - 14)(3D^6 - 108D^5 + 1586D^4 - 11304D^3 + 41928D^2 - 78208D + 57984)}{10(5D - 16)(5D - 18)(5D - 22)(D - 3)} \right) \\
&+ A_{9,4} \left( + \frac{1701D^3}{160000} - \frac{375273D^2}{1600000} + \frac{5426811D}{3200000} + \frac{4067}{1536(2D - 7)} + \frac{232399}{665600(2D - 9)} \right) \\
&\quad - \frac{8}{25(D - 2)} - \frac{675}{896(D - 3)} + \frac{4194344}{1015625(5D - 16)} - \frac{287008}{78125(5D - 18)} - \frac{3778904}{1640625(5D - 22)}
\end{align*}
\]
\[- \frac{116599473}{32000000} \]

\[X_{C_f C_A}^q = \]

\[-B_{4,1} \left( + \frac{153701 D^3}{625} - \frac{45111262 D^2}{9375} + \frac{3307905503 D}{112500} - \frac{7045}{165455} + \frac{140183197}{44164} \right) \]

\[- \frac{27(3D - 10)}{2060} + \frac{55(3D - 14)}{9383166} - \frac{2673(D - 1)}{2673} + \frac{25(D - 2)}{25} + \frac{12(D - 3)}{12} + \frac{3748279}{143} \]

\[+ \frac{4860(D - 4)}{13826309401} - \frac{402187500(5D - 16)}{8793673} - \frac{234375(5D - 18)}{915068} - \frac{20646}{345128} - \frac{128}{20864} \]

\[+ \frac{12(D - 3)^2}{31855488829} + \frac{405(D - 4)^2}{12} + \frac{135(D - 4)^3}{45} + \frac{45(D - 4)^4}{5(D - 4)^5} \]

\[- \frac{843750}{3} \]

\[-A_{5,1} \left( + \frac{4277 D^3}{125} - \frac{243647 D^2}{375} + \frac{106547887 D}{28125} - \frac{49315}{24} + \frac{18960447}{2600(2D - 9)} \right) \]

\[- \frac{1323(3D - 10)}{357584} + \frac{567567(D - 1)}{7618840} + \frac{25(D - 2)}{10064} + \frac{3(D - 3)}{10} + \frac{405(D - 4)}{370389} \]

\[+ \frac{1340625(5D - 16)}{563608248} + \frac{203125(5D - 18)}{1415791832} - \frac{46875(5D - 22)}{189(D - 1)^2} + \frac{81212}{45(D - 4)^2} + \frac{11584}{843750} \]

\[+ \frac{45(D - 4)^3}{256} - \frac{3(D - 4)^4}{67500} \]

\[+ A_{5,2} \left( + \frac{20594 D^3}{1875} - \frac{6630308 D^2}{28125} + \frac{25700042 D}{16875} - \frac{1409}{6(2D - 7)} + \frac{2348211}{650(2D - 9)} \right) \]

\[- \frac{567(3D - 10)}{11944} + \frac{99(3D - 14)}{801980} - \frac{243243(D - 1)}{10360820} + \frac{225(D - 2)}{75128} - \frac{1215(D - 4)}{559813} \]

\[+ \frac{33515625(5D - 16)}{1052157372} - \frac{1015625(5D - 18)}{1170515298} - \frac{703125(5D - 22)}{64} - \frac{306428}{(D - 2)^2} - \frac{3344023858}{405(D - 4)^2} \]

\[- \frac{135(D - 4)^3}{1408} - \frac{15(D - 4)^4}{1265625} \]

\[-B_{5,1} \left( + \frac{7497 D^3}{1250} - \frac{394839 D^2}{3125} + \frac{10983001 D}{12500} - \frac{75999}{712201} + \frac{1424}{200430972} - \frac{146129984}{78125(5D - 18)} \right) \]

\[- \frac{1}{(D - 3)} + \frac{1055}{(D - 4)} - \frac{234375(5D - 16)}{78125(5D - 18)} - \frac{189456643}{125000} \]

\[+ \frac{5(D - 4)^2}{4176} + \frac{296}{(D - 4)} + \frac{1152}{5(D - 4)^4} \]

\[+ B_{5,2} \left( + \frac{13666 D^3}{625} - \frac{4406452 D^2}{9375} + \frac{92206756 D}{28125} - \frac{13818}{3748279} + \frac{320}{169244232} + \frac{122152296}{27(3D - 10)} \right) \]

\[- \frac{2536}{6(D - 4)} + \frac{10325}{15625(5D - 16)} + \frac{169244232}{78125(5D - 18)} + \frac{78125(5D - 22)}{78125(5D - 22)} \]

\[+ \frac{22774}{15(D - 4)^2} + \frac{15736}{15(D - 4)^3} + \frac{7552}{15(D - 4)^4} - \frac{2504444962}{421875} \]
\[
\begin{align*}
&= -\frac{38}{45(D-2)} + \frac{1833}{140(D-3)} + \frac{8}{5(D-4)} - \frac{732913468}{9140625(5D-16)} - \frac{368278}{234375(5D-18)} + \frac{4921875(5D-22)}{71838976} - \frac{1}{12(D-3)^2} + \frac{1}{16428169} \\
&- A_{8,1} \left( + \frac{1107D^3}{160000} + \frac{110409D^2}{1600000} - \frac{2547331D}{640000} + \frac{2273}{899587} + \frac{56}{234375(5D-16)} - \frac{71838976}{78125(5D-18)} + \frac{493416}{546875(5D-22)} - \frac{9863}{1024(2D-7)^2} + \frac{1}{3200000} \right) \\
&+ A_{9,1} \left( + \frac{243D^3}{5000} - \frac{62289D^2}{50000} + \frac{283527D}{25000} + \frac{88}{5(D-2)} - \frac{285}{56(D-3)} + \frac{2}{(D-4)} + \frac{1549163}{38866491} + \frac{1}{3200000} \right) \\
&- A_{9,2} \left( + \frac{567D^3}{160000} - \frac{125091D^2}{1600000} + \frac{180893D}{320000} + \frac{4067}{3046875(5D-16)} - \frac{287008}{234375(5D-18)} + \frac{3778904}{4921875(5D-22)} - \frac{38866491}{3200000} \right) \\
&- A_{9,4} \left( + \frac{8}{75(D-2)} - \frac{225}{896(D-3)} + \frac{4194344}{3046875(5D-16)} - \frac{287008}{234375(5D-18)} + \frac{3778904}{4921875(5D-22)} - \frac{38866491}{3200000} \right)
\end{align*}
\]

\[
X_{C_2^pN_F}^q = \begin{align*}
&+ B_{4,1} \left( + \frac{72D^2}{5} - \frac{11524D}{75} - \frac{37120}{63(3D-10)} - \frac{742964}{6237(D-1)} + \frac{4}{(D-3)} + \frac{10576}{81(D-4)} + \frac{319872}{1375(5D-16)} - \frac{3088}{27(D-4)^2} + \frac{1024}{9(D-4)^3} + \frac{256}{3(D-4)^4} + \frac{412948}{1125} \right) \\
&- A_{5,1} \left( + \frac{1216D^2}{75} - \frac{83936D}{375} - \frac{512}{3(3D-10)} - \frac{2293120}{11583(D-1)} - \frac{19648}{81(D-4)} + \frac{297024}{6875(5D-16)} + \frac{3698688}{8125(5D-18)} - \frac{6272}{27(D-4)^2} + \frac{1024}{9(D-4)^3} + \frac{983968}{1875} \right) \\
&- A_{5,2} \left( + \frac{4D^2}{75} - \frac{78712D}{1125} - \frac{6656}{189(3D-10)} - \frac{32}{9(3D-8)} - \frac{4357048}{81081(D-1)} + \frac{20464}{131376} + \frac{1585152}{1585152} - \frac{5312}{27(D-4)^2} - \frac{1024}{9(D-4)^3} + \frac{2009608}{16875} \right)
\end{align*}
\]
\[-A_{6,1} \frac{(D^2 - 7D + 16)(6D^3 - 65D^2 + 238D - 288)(D - 2)}{2(D - 3)(D - 1)(D - 4)^2} + A_{6,3} \left( \frac{28D^2}{25} - \frac{2868D}{125} - \frac{32}{45(3D - 8)} - \frac{500924}{19305(D - 1)} - \frac{400}{27(D - 4)} - \frac{39984}{6875(5D - 16)} - \frac{176128}{8125(5D - 18)} - \frac{128}{9(D - 4)^2 + 334516} \right) \]
\[-B_{6,2} \frac{(D^2 - 7D + 16)(3D^3 - 31D^2 + 110D - 128)(D - 2)}{(D - 1)(3D - 8)(D - 4)^2} - C_{6,1} \frac{(D^2 - 7D + 16)(3D^3 - 31D^2 + 110D - 128)(D - 2)}{(D - 1)(3D - 8)(D - 4)^2} - A_{7,2} \frac{8(D - 2)(D^4 - 28D^3 + 220D^2 - 696D + 784)}{(5D - 18)(D - 1)(5D - 16)} \]

\[
X_{CFCA\bar{N}F}^2 =
B_{4,1} \left( + \frac{24D^2}{5} - \frac{396D}{25} - \frac{250}{3(3D - 10)} + \frac{330910}{2673(D - 1)} - \frac{170}{3(D - 2)} + \frac{34408}{92064} + \frac{1375(5D - 16)}{3(D - 2)^2} + \frac{32}{14624} + \frac{256}{27(D - 4)^3} \right)
+ A_{5,1} \left( + \frac{208D^2}{75} - \frac{3656D}{125} + \frac{25352}{441(3D - 10)} - \frac{4990592}{63063(D - 1)} + \frac{184}{3(D - 2)} - \frac{992}{148512} + \frac{8125(5D - 18)}{189(D - 1)^2} - \frac{6976}{(D - 2)^2} \right)
+ A_{5,2} \left( + \frac{84D^2}{25} - \frac{6984D}{125} + \frac{460}{63(3D - 10)} - \frac{20721640}{243243(D - 1)} + \frac{20}{3(D - 2)} - \frac{24568}{24312} + \frac{8125(5D - 18)}{792576} - \frac{16}{5024} \right)
+ A_{6,1} \left( + \frac{3D^2}{2} - \frac{47D}{4} - \frac{3073}{72(D - 1)} + \frac{86}{3(D - 2)} + \frac{11}{8(D - 3)} + \frac{8}{9(D - 4)} - \frac{109}{4(D - 1)^2} - \frac{8}{(D - 2)^2} + \frac{16}{3(D - 4)^2} + \frac{133}{4} \right)
+ A_{6,2} \frac{90D^7 - 1803D^6 + 15301D^5 - 70848D^4 + 191676D^3 - 299024D^2 + 242976D - 74880}{9(5D - 16)(D - 3)(D - 2)^2(D - 1)(D - 4)}
- A_{6,3} \frac{42D^7 - 656D^6 + 3854D^5 - 11430D^4 + 24896D^3 - 65144D^2 + 134560D - 113856}{3(D - 1)(D - 2)^2(5D - 18)(5D - 16)}
\]
\[ +A_{7,2} \frac{4(D - 2)(D^4 - 28D^3 + 220D^2 - 696D + 784)}{(5D - 18)(D - 1)(5D - 16)} \]

\[ X_{C_F N_F}^q = +A_{6,1} \frac{(6D^3 - 65D^2 + 238D - 288)(D - 2)^2}{6(D - 4)(D - 3)(D - 1)^2} \]

\[ X_{C_F N_{F,V}}^q = +B_{4,1} \left( \frac{119132D}{625} + \frac{380}{9(2D - 7)} + \frac{50245888}{975(2D - 9)} + \frac{280}{3(3D - 10)} + \frac{9088443}{242(3D - 14)} \right) + B_{4,1} \left( \frac{9269061200}{722007(D - 1)} - \frac{4169543}{675(D - 2)} + \frac{182}{3(D - 3)} + \frac{1194157}{54(D - 4)} - \frac{57818921783}{265443750(5D - 16)} \right) + B_{4,1} \left( \frac{80888593}{243750(5D - 18)} + \frac{53125(5D - 22)}{221659776} + \frac{9(D - 2)^2}{15748} + \frac{35}{3(D - 3)^2} + \frac{750554}{45(D - 4)^2} - \frac{3(3D - 2)^3}{608 + 25148 + 2048 + \frac{12730856}{9375}} \right) + A_{5,1} \left( -\frac{12243D}{625} + \frac{665}{9(2D - 7)} - \frac{468624}{325(2D - 9)} - \frac{56}{3(3D - 10)} - \frac{92219200}{65637(D - 1)} \right) - A_{5,2} \left( -\frac{146444D}{1875} - \frac{76}{9(2D - 7)} + \frac{40784}{325(2D - 9)} + \frac{56}{9(3D - 10)} + \frac{215775}{121(3D - 14)} \right) + B_{5,1} \left( \frac{3474D}{625} - \frac{196273}{325(2D - 9)} - \frac{25123840}{65637(D - 1)} + \frac{144542}{225(D - 2)} - \frac{2}{(D - 3)} + \frac{135(D - 4)}{93338} + \frac{1437614}{14095976} + \frac{190269}{39796848} + \frac{1984}{3(D - 2)^2} + \frac{2900}{135(D - 4)^2} + \frac{224}{(D - 2)^3} + \frac{1664}{15(D - 4)^3} + \frac{190269}{3125} \right) - B_{5,2} \left( \frac{37136D}{1875} - \frac{285488}{325(2D - 9)} + \frac{80}{9(3D - 10)} + \frac{28564480}{65637(D - 1)} + \frac{163559}{450(D - 2)} \right) \]
\[
\begin{align*}
- \frac{150403}{135(D - 4)} &+ \frac{2567823}{4021875(5D - 16)} - \frac{13477879}{81250(5D - 18)} + \frac{10007466}{53125(5D - 22)} - \frac{2972}{3(D - 2)^2} \\
+ \frac{32612}{45(D - 4)^2} &+ \frac{14385}{(D - 2)^3} - \frac{58221088}{18(D - 2)} + \frac{1}{2(D - 3)} \\
- \frac{480}{418} &+ \frac{3328}{3759722} + \frac{4577488}{432134} + \frac{28125}{3537842} \\
\frac{135(D - 4)}{64} &+ \frac{3403125(5D - 16)}{40625(5D - 18)} + \frac{53125(5D - 22)}{(D - 2)^2} \\
\frac{64}{18750} &+ \frac{100399}{100399} \\
- \frac{17533D}{3000} &+ \frac{9}{24(D - 7)} + \frac{10196}{975(2D - 9)} + \frac{1382570}{1683(D - 1)} - \frac{411274}{675(D - 2)} \\
- \frac{1623}{56(D - 3)} &+ \frac{1088}{45(D - 4)} + \frac{18324568}{2413125(5D - 16)} + \frac{1024224}{74375(5D - 22)} + \frac{3680}{9(D - 2)^2} \\
- \frac{11}{6(D - 3)^2} &+ \frac{256}{3(D - 2)^3} + \frac{428883}{1875} \\
+ \frac{26253D}{2500} &+ \frac{95}{36(D - 7)} - \frac{41793584}{21879(D - 1)} - \frac{13967}{9(D - 2)} + \frac{3}{(D - 3)} \\
+ \frac{3548}{45(D - 4)} &+ \frac{39662}{103125(5D - 16)} + \frac{53125(5D - 22)}{883413} + \frac{12129744}{3(D - 2)^2} + \frac{2048}{12500} \\
+ \frac{352}{45(D - 4)^2} &+ \frac{243387}{243387} \\
+ \frac{15(D - 4)^2}{(D - 2)^3} &+ \frac{12500}{3(D - 2)^2} \\
+ \frac{9D}{100} &+ \frac{37}{18(2D - 7)} + \frac{2549}{1560(2D - 9)} + \frac{143(D - 1)}{143(D - 1)} + \frac{280}{90(D - 2)} \\
\frac{144076}{160875(5D - 16)} &+ \frac{6313}{9750(5D - 18)} + \frac{4}{(D - 2)^2} + \frac{3927}{1000} \\
- \frac{5D^6 - 93D^5 + 892D^4 - 4656D^3 + 12528D^2 - 15472D + 6080}{(D - 1)(D - 2)^2(5D - 18)(5D - 16)} \\
- \frac{A_{7,1}}{2} &+ \frac{211D}{625} - \frac{37}{18(2D - 7)} + \frac{2549}{5850(2D - 9)} - \frac{5424}{187(D - 1)} + \frac{2566}{75(D - 2)} \\
+ \frac{3(D - 3)}{2} &+ \frac{5850(2D - 9)}{26995456} - \frac{79688}{79688} + \frac{56582}{56582} - \frac{328}{3(D - 2)^2} \\
- \frac{128}{(D - 3)^2} &+ \frac{4576}{3125} \\
+ \frac{A_{7,2}}{3(D - 3)^2} &+ \frac{3021875(5D - 16)}{28125(5D - 18)} - \frac{53125(5D - 22)}{53125(5D - 22)} - \frac{3(D - 2)^2}{3(D - 2)^2} \\
+ \frac{8}{12(D - 3)^2} &+ \frac{273763}{75000} \\
+ \frac{A_{7,3}}{12(D - 3)^2} &+ \frac{162D}{625} - \frac{37}{72(2D - 7)} + \frac{2549}{1170(2D - 9)} - \frac{22736}{2431(D - 1)} + \frac{754}{45(D - 2)} \\
- \frac{323}{84(D - 3)} &+ \frac{2696828}{2696828} - \frac{97442}{97442} + \frac{1877744}{1877744} - \frac{8}{12(D - 3)^2} \\
+ \frac{1}{12(D - 3)^2} &+ \frac{273763}{75000} \\
+ \frac{A_{7,4}}{12(D - 3)^2} &+ \frac{162D}{625} - \frac{37}{72(2D - 7)} + \frac{2549}{1170(2D - 9)} - \frac{22736}{2431(D - 1)} + \frac{754}{45(D - 2)}
\end{align*}
\]
$$\frac{-323}{84(D - 3)} - \frac{269628}{4021875(5D - 16)} + \frac{97442}{365625(5D - 18)} - \frac{1877744}{1115625(5D - 22)} - \frac{8}{(D - 2)^2}$$

$$- \frac{1}{12(D - 3)^2} + \frac{273763}{75000}$$

$$+ A_{8,1} \left( + \frac{1107D}{80000} + \frac{1643}{22363} - \frac{8085}{77484} + \frac{349}{85352} + \frac{1305}{133} \right.$$  

$$\left. + \frac{28125(5D - 16)}{1511991} - \frac{40625(5D - 18)}{80000} \right)$$

$$- A_{9,1} \left( + \frac{243D}{5000} + \frac{24255}{83136} - \frac{24}{56(D - 3)} - \frac{75}{(D - 2)} - \frac{10076}{9375(5D - 16)} \right.$$  

$$\left. - \frac{40625(5D - 18)}{80000} + \frac{252703}{1115625(5D - 22)} + \frac{16}{5103} \right)$$

$$+ A_{9,2} \left( \frac{3(3D - 14)(D - 4)(D^5 + 25D^4 - 290D^3 + 1036D^2 - 1560D + 928)}{20(D - 1)(D - 2)(D - 3)(5D - 22)(5D - 18)(5D - 16)} \right)$$

$$X_9^{\sigma_{C_4}} =$$

$$+ B_{4,1} \left( + \frac{282D^2}{125} - \frac{2592693D}{106048} - \frac{4480}{63(2D - 7)} - \frac{447830}{37565622628} + \frac{46485469184}{16575(2D - 9)} \right.$$  

$$\left. + \frac{21(3D - 10)}{1120} - \frac{55(3D - 14)}{635092036} + \frac{5(3D - 8)}{1440} - \frac{1216215(D - 1)}{29989344832} + \frac{225(D - 2)}{5361643396} \right.$$  

$$\left. + \frac{3(D - 3)}{6939682834} - \frac{1215(D - 4)}{10119933824} + \frac{5(D - 6)}{55039} - \frac{30121875(D - 14)}{102944} + \frac{4021875(5D - 16)}{1787} \right.$$  

$$\left. - \frac{121875(5D - 18)}{130962128} + \frac{3125(5D - 22)}{3328} + \frac{108(D - 1)^2}{21706624} - \frac{5(D - 2)^2}{359296} - \frac{5(D - 3)^2}{2048} \right.$$  

$$\left. - \frac{405(D - 4)^2}{600552286} + \frac{(D - 2)^3}{9375} - \frac{135(D - 4)^3}{3(D - 4)^3} - \frac{9(D - 4)^4}{3(D - 4)^4} + \frac{3(D - 4)^5}{9375} \right)$$

$$- A_{5,1} \left( + \frac{2114656D}{375} + \frac{608}{5(2D - 5)} + \frac{13520}{9(2D - 7)} - \frac{415048832}{5525(2D - 9)} - \frac{256192}{2205(3D - 10)} \right.$$  

$$\left. - \frac{6487807808}{567567(2D - 1)} + \frac{8756344}{225(D - 2)} + \frac{32}{(D - 3)} + \frac{5951984}{81(D - 4)} + \frac{5760}{7(D - 6)} \right.$$  

$$\left. + \frac{4303125(5D - 14)}{262210816} + \frac{7942807088}{3072} + \frac{776706728}{1263488} - \frac{561369536}{23552} - \frac{22016}{189(D - 1)^2} \right.$$  

$$\left. - \frac{5(D - 2)^2}{993704848} + \frac{3(D - 4)^2}{28125} + \frac{(D - 2)^3}{3(D - 4)^3} + \frac{45(D - 4)^3}{3(D - 4)^4} + \frac{3(D - 4)^4}{993704848} \right)$$

$$- A_{5,2} \left( + \frac{1594502D}{1125} - \frac{576}{1485(2D - 5)} + \frac{37696}{63(2D - 7)} - \frac{10816}{5525(2D - 9)} - \frac{37731712}{1125} \right.$$  

$$- \frac{1125}{120(2D - 3)^2} + \frac{273763}{75000}$$
\[-B_{5,1} \left( +204D^2 - \frac{495373D}{250} + \frac{5681}{72(2D - 5)} + \frac{13125}{8(2D - 7)} - \frac{18158364}{5525(2D - 9)} \right) \]
\[-\frac{136352}{945(3D - 10)} - \frac{101318}{33(3D - 14)} + \frac{448}{15(3D - 8)} - \frac{1468089542}{405405(3D - 1)} + \frac{4190624}{225(3D - 2)} + \frac{8039792}{1215(3D - 4)} + \frac{994}{7(3D - 6)} + \frac{12736}{90365625(3D - 14)} + \frac{28153125(5D - 16)}{28153125(5D - 16)} + \frac{3125(5D - 22)}{3125(5D - 22)} - \frac{27(D - 1)^2}{141384312} - \frac{(D - 2)^2}{(D - 2)^2} + \frac{405(5D - 4)^2}{405(5D - 4)^2} + \frac{(D - 2)^3}{(D - 2)^3} + \frac{135(5D - 4)^3}{135(5D - 4)^3} + \frac{28160}{902670584} - \frac{9(D - 4)^3}{9(D - 4)^3} - \frac{84375}{84375} \]

\[+ B_{5,2} \left( + \frac{437838D}{125} + \frac{3078304}{10395(5D - 5)} - \frac{264121984}{5525(2D - 9)} - \frac{19136}{105(3D - 10)} + \frac{128}{5(3D - 8)} \right) \]
\[-\frac{3842871566}{100409312} + \frac{2331716}{75(5D - 2)} + \frac{2263616}{135(5D - 4)} + \frac{288}{7(5D - 6)} - \frac{405405(3D - 1)}{405405(3D - 1)} + \frac{1340625(5D - 16)}{40625(5D - 16)} - \frac{3125(5D - 22)}{3125(5D - 18)} - \frac{9375(5D - 22)}{9375(5D - 22)} - \frac{5(5D - 2)^2}{5(5D - 2)^2} + \frac{5(5D - 4)^2}{5(5D - 4)^2} + \frac{(D - 2)^3}{(D - 2)^3} + \frac{315(5D - 4)^3}{315(5D - 4)^3} + \frac{7680}{146133933} - \frac{3125}{3125} \]

\[+ A_{6,1} \left( +D^2 + \frac{6324D}{125} + \frac{101318}{99(3D - 14)} + \frac{1186}{81(5D - 1)} - \frac{716}{3(3D - 2)} \right) \]
\[-\frac{796}{3(3D - 4)} + \frac{6706784}{253125(5D - 14)} - \frac{623884}{34375(5D - 16)} - \frac{38922}{3125(5D - 18)} - \frac{5049408}{3125(5D - 22)} + \frac{9684408}{9(3D - 1)} - \frac{112372D}{225} - \frac{640}{21(2D - 5)} + \frac{338}{21(2D - 7)} + \frac{9432928}{16575(2D - 9)} \]
\[-\frac{103424}{5(2D - 1)} + \frac{365056}{75(5D - 2)} + \frac{16544}{105(3D - 3)} - \frac{23576}{135(5D - 4)} - \frac{343264}{7(5D - 6)} - \frac{3328}{3843264} - \frac{40}{(D - 2)^2} - \frac{40}{(D - 3)^2} + \frac{14730632}{14730632} - \frac{5625}{5625} \]

\[+ A_{6,2} \left( +4D^2 + \frac{133488D}{125} - \frac{3380}{63(2D - 7)} + \frac{1216}{15(3D - 8)} + \frac{9376286}{19305(3D - 1)} + \frac{525424}{125} + \frac{97276}{50} - \frac{118752128}{11452108} + \frac{11452108}{309375(5D - 16)} \right) \]
\[
\begin{align*}
&\frac{77648}{40625(5D - 18)} - \frac{17312256}{3125(5D - 22)} - \frac{42752}{5(D - 2)^2} + \frac{12496}{15(D - 4)^2} + \frac{1536}{(D - 2)^3} + \\
&\frac{128}{(D - 4)^3} - \frac{74684696}{9375} \right)
-B_{6,1} \left( +4D^3 - 120D^2 + 1200D + \frac{2112}{(D - 2)} - \frac{768}{(D - 4)} \right) \\
&- \frac{1152}{(D - 2)^2} - \frac{768}{(D - 4)^2} + \frac{256}{(D - 2)^3} - \frac{256}{(D - 4)^3} - 4096 \right)
-B_{6,2} \left( +54D^2 - \frac{1689D}{2} + \frac{2261}{72(2D - 5)} + \frac{625}{8(2D - 7)} + \frac{1400}{9(D - 1)} \right) \\
&- \frac{3152}{(D - 2)} - \frac{2440}{3(D - 4)} + \frac{2048}{(D - 2)^2} - \frac{2800}{3(D - 4)^2} - \frac{512}{(D - 2)^3} \\
&- \frac{384}{(D - 4)^3} + 3530 \right)
-C_{6,1} \left( +54D^2 - \frac{1689D}{2} + \frac{2261}{72(2D - 5)} + \frac{625}{8(2D - 7)} - \frac{1400}{9(D - 1)} \right) \\
&- \frac{3152}{(D - 2)} - \frac{2440}{3(D - 4)} + \frac{2048}{(D - 2)^2} - \frac{2800}{3(D - 4)^2} - \frac{512}{(D - 2)^3} \\
&- \frac{384}{(D - 4)^3} + 3530 \right)
-A_{7,1} \left( +\frac{54D}{5} + \frac{304}{77(2D - 5)} - \frac{4792}{63(2D - 7)} + \frac{294779}{3315(2D - 9)} - \frac{9072}{715(D - 1)} \right) \\
&- \frac{1396}{45(D - 2)} + \frac{80}{(D - 4)} + \frac{48}{35(D - 6)} + \frac{228}{119(5D - 14)} - \frac{32175}{32175(5D - 16)} \\
&+ \frac{557732}{10725(5D - 18)} + \frac{32}{(D - 2)^2} + \frac{873}{25} \right)
+A_{7,2} \left( +\frac{784D}{25} - \frac{12392}{715(5D - 1)} + \frac{104}{(D - 2)} - \frac{160}{(D - 4)} - \frac{192}{35(D - 6)} \right) \\
&+ \frac{704}{125(5D - 14)} + \frac{21216}{1925(5D - 16)} + \frac{488216}{1625(5D - 18)} - \frac{64}{(D - 2)^2} \right)
+A_{7,3} \left( +\frac{26064D}{125} + \frac{228}{(2D - 5)} - \frac{4792}{63(2D - 7)} + \frac{1179116}{49725(2D - 9)} - \frac{2008}{(D - 1)} \right) \\
&+ \frac{1444072}{225(D - 2)} + \frac{17032}{35(D - 3)} - \frac{56929472}{74375(5D - 14)} - \frac{2013568}{73125(5D - 16)} + \frac{707192}{5625(5D - 18)} \\
&- \frac{656704}{13125(5D - 22)} - \frac{7360}{(D - 2)^2} - \frac{40}{(D - 3)^2} + \frac{2048}{(D - 2)^3} - \frac{1315312}{625} \right)
-A_{7,4} \left( +\frac{8401D}{250} - \frac{19}{4(2D - 5)} + \frac{1198}{63(2D - 7)} + \frac{9945}{9945(2D - 9)} + \frac{43}{15(D - 1)} \right) \\
&+ \frac{4936}{45(D - 2)} - \frac{717}{35(D - 3)} + \frac{16}{5(D - 6)} + \frac{44793584}{1115625(5D - 14)} - \frac{4939984}{365625(5D - 16)} \\
&- \frac{478744}{28125(5D - 18)} - \frac{7045984}{65625(5D - 22)} - \frac{64}{(D - 2)^2} + \frac{1}{(D - 3)^2} - \frac{1219407}{12500} \right)
\]
\[
\begin{align*}
& \frac{44032}{189(D - 1)^2} - \frac{19712}{3(D - 2)^2} - \frac{250816}{135(D - 4)^2} + \frac{1024}{(D - 2)^3} - \frac{77056}{45(D - 4)^3} \\
& + \frac{16875}{183292489}
\end{align*}
\]

\[+ A_{5,2} \left( + \frac{30752D^2}{375} - \frac{3289648D}{1875} + \frac{384}{D} + \frac{150784}{1485(2D - 5)} - \frac{1720}{63(2D - 7)} \\
- \frac{95525(2D - 9)}{444172} - \frac{945(3D - 10)}{1509632} + \frac{231(3D - 14)}{3976} + \frac{135(3D - 8)}{269380672} + \frac{135135(5D - 1)}{29517696}
\right)
\]

\[+ B_{5,1} \left( + \frac{1397D^2}{250} + \frac{64933D}{2500} - \frac{561}{18(2D - 5)} - \frac{6125}{24(2D - 7)} - \frac{8679748}{5525(2D - 9)} \\
+ \frac{1269580}{1954816} - \frac{225(2D - 2)}{2} - \frac{45(5D - 4)}{4984} + \frac{4303125(5D - 14)}{57180032} + \frac{155(5D - 18)}{12544}
\right)
\]

\[+ B_{5,2} \left( + \frac{3484D^2}{75} - \frac{6867848D}{5625} + \frac{12313216}{10395(2D - 5)} + \frac{12625088}{5525(2D - 9)} - \frac{1024}{945(3D - 10)} \\
- \frac{45(3D - 8)}{6733664} + \frac{405405(5D - 1)}{551540128} + \frac{225(2D - 2)}{177548474} + \frac{405(5D - 4)}{27136592} - \frac{7(D - 6)}{21392}
\right)
\]

\[+ A_{6,1} \left( + \frac{394D^2}{25} - \frac{5986D}{25} - \frac{26416}{693(3D - 14)} + \frac{17332}{81(D - 1)} - \frac{636}{5(D - 2)} \\
+ \frac{15400288}{176} - \frac{354375(5D - 14)}{4125(5D - 16)} + \frac{81(D - 1)}{18256} - \frac{46252}{625(5D - 18)} + \frac{70784}{1875(5D - 22)}
\right)
\]

\[+ A_{6,2} \left( + \frac{238D^2}{25} - \frac{84531D}{250} - \frac{1280}{3D} - \frac{4408}{21(2D - 5)} + \frac{215}{84(2D - 7)} \\
- \frac{93023}{1950(2D - 9)} - \frac{206848}{297(D - 1)} - \frac{77056}{75(D - 2)} + \frac{105(D - 3)}{8366} - \frac{135(D - 4)}{1324688}
\right)
\]

\[+ \frac{1088}{7(D - 6)} + \frac{118125(5D - 14)}{5630625(5D - 16)} - \frac{645725648}{39375(5D - 22)} + \frac{1088}{3(D - 2)^2}
\]
\[-A_{6,3} \left( + \frac{16}{(D - 3)^2} - \frac{256}{3(D - 2)^3} + \frac{40336667}{22500} \right) \]
\[-B_{6,2} \left( + \frac{7D^2}{2} - \frac{185D}{4} + \frac{2261}{18(2D - 5)} + \frac{875}{72(2D - 7)} - \frac{2800}{9(D - 1)} \right) \]
\[-C_{6,1} \left( + \frac{7D^2}{2} - \frac{185D}{4} + \frac{2261}{18(2D - 5)} + \frac{875}{72(2D - 7)} - \frac{2800}{9(D - 1)} \right) \]
\[-A_{7,1} \left( + \frac{16D^2}{25} - \frac{1437D}{125} + \frac{1216}{77(2D - 5)} - \frac{902}{63(2D - 7)} - \frac{28181}{9012656} + \frac{152185504}{2432} \right) \]
\[+A_{7,2} \left( + \frac{17D^2}{20} - \frac{2237D}{100} + \frac{71585}{2016(2D - 7)} - \frac{226533}{35360(2D - 9)} - \frac{420416}{715(2D - 1)} \right) \]
\[+A_{7,3} \left( + \frac{44D^2}{125} - \frac{19864D}{625} + \frac{912}{(2D - 5)} - \frac{902}{63(2D - 7)} - \frac{56362}{49725(2D - 9)} \right) \]
\[-\frac{4016}{(D - 1)} + \frac{566216}{225(2D - 2)} - \frac{808}{5(D - 3)} + \frac{252301792}{371875(5D - 14)} + \frac{6857728}{73125(5D - 16)} \]
\[+A_{7,4} \left( + \frac{21D^2}{250} - \frac{15043D}{1250} - \frac{19}{(2D - 5)} + \frac{6805}{1008(2D - 7)} - \frac{93023}{9360(2D - 9)} \right) \]
\[+ \frac{86}{15(D - 1)} - \frac{3116}{45(D - 2)} + \frac{934}{105(D - 3)} + \frac{8}{5(D - 6)} + \frac{4403312}{65625(5D - 14)} \]
\[
- \frac{10487648}{365625(5D - 16)} - \frac{17612}{5625(5D - 18)} + \frac{1574056}{196875(5D - 22)} + \frac{1}{(D - 3)^2} + \frac{916069}{15000}
\]

\[
-A_{7,5} \left( + \frac{21D^2}{250} - \frac{15043D}{1250} - \frac{19}{2(5D - 5)} + \frac{1908(2D - 7)}{9360(2D - 9)} - \frac{1}{(D - 3)^2} + \frac{93023}{4403312} \right)
\]

\[
+ \frac{15(D - 1)}{3116} - \frac{45(D - 2)}{934} + \frac{105(D - 3)}{8} + \frac{5(D - 6)}{65625(5D - 14)}
\]

\[
- \frac{10487648}{365625(5D - 16)} - \frac{17612}{5625(5D - 18)} + \frac{1574056}{196875(5D - 22)} + \frac{1}{(D - 3)^2} + \frac{916069}{15000}
\]

\[
-A_{8,1} \left( + \frac{113787}{20000} - \frac{2377}{448(2D - 7)} - \frac{2848}{45(2D - 2)} - \frac{7(D - 3)}{21875(5D - 14)} - \frac{1}{(D - 3)^2} + \frac{679432}{20000} \right)
\]

\[
+ \frac{28125(5D - 16)}{51216} + \frac{3125(5D - 18)}{65625(5D - 22)} - \frac{192(2D - 7)^2}{920019}
\]

\[
-B_{8,1} \frac{(2D^3 - 25D^2 + 94D - 112)(D^3 - 16D^2 + 68D - 88)}{2(2D - 5)(2D - 7)(D - 2)^2}
\]

\[
-C_{8,1} \frac{(2D^3 - 25D^2 + 94D - 112)(D^3 - 16D^2 + 68D - 88)}{2(2D - 5)(2D - 7)(D - 2)^2}
\]

\[
+A_{9,1} \left( + \frac{81D^2}{250} - \frac{9369D}{1250} + \frac{176}{3(D - 2)} + \frac{45}{14(D - 3)} + \frac{32928}{625(5D - 14)} \right)
\]

\[
- \frac{160688}{9375(5D - 16)} + \frac{3125(5D - 18)}{65625(5D - 22)} + \frac{12424}{1250} + \frac{54801}{125}
\]

\[
+A_{9,2} \frac{(3D - 14)(D - 4)(135D^5 - 2200D^4 + 15156D^3 - 54336D^2 + 99776D - 74112)}{5(5D - 3)(5D - 22)(5D - 16)(5D - 14)(5D - 18)(D - 2)}
\]

\[
-A_{9,4} \left( + \frac{81D^2}{4000} - \frac{513D}{2500} - \frac{973}{2304(2D - 7)} - \frac{75511}{1414400(2D - 9)} - \frac{16}{75(D - 2)} \right)
\]

\[
+ \frac{3024}{10625(5D - 14)} - \frac{56144}{121875(5D - 16)} + \frac{4672}{3125(5D - 18)} + \frac{8936}{28125(5D - 22)} + \frac{36117}{80000}
\]

\[
X^9_{C_{A}C_{F}N_{P}} =
\]

\[
+B_{4,1} \left( + \frac{676832D^2}{125} - \frac{114134744D}{1875} - \frac{4480}{3D} + \frac{17860}{7(2D - 7)} + \frac{4283546344}{16575(2D - 9)} \right)
\]

\[
- \frac{226592}{63(3D - 10)} + \frac{36336222}{385(3D - 14)} - \frac{6656}{45(3D - 8)} + \frac{79424288}{405405(D - 1)} - \frac{13370848}{675(D - 2)}
\]

\[
- \frac{4552}{3(3D - 3)^2} + \frac{1922612}{135(D - 4)} + \frac{8647206224}{30121875(5D - 14)} + \frac{3472406764}{2413125(5D - 16)} + \frac{1338317458}{24375(5D - 18)}
\]

\[
- \frac{6431698288}{9375(5D - 22)} + \frac{53504}{9(2D - 2)^2} - \frac{76}{3(D - 3)^2} + \frac{4864}{15(D - 4)^2} + \frac{9560}{15(D - 4)^3}
\]

\[
+ \frac{187136}{5(5D - 4)^4} + \frac{3507215464}{28125}
\]

\[
+A_{5,1} \left( + \frac{103212D^2}{125} - \frac{43508066D}{5625} + \frac{4864}{5(2D - 5)} + \frac{9560}{9(2D - 7)} + \frac{64410511}{5525(2D - 9)} \right)
\]
\[
\begin{align*}
A_{5.2} &= -\frac{255904}{45502976} - \frac{55040}{26533088} + \frac{1180528}{38430896} + \frac{160}{3(D - 3)} - \frac{2847392}{28747264} - \frac{14848}{3(D - 2)^2} \\
B_{5.1} &= +\frac{109888D^2}{375} - \frac{27257176D}{5625} - \frac{192}{D} + \frac{301568}{1485(2D - 5)} - \frac{7648}{63(2D - 7)} \\
B_{5.2} &= -\frac{29636D^2}{75} - \frac{21234136D}{5625} + \frac{24626432}{10395(2D - 5)} + \frac{1571584}{325(2D - 9)} + \frac{114112}{945(3D - 10)} \\
A_{6.1} &= +\frac{21284D^2}{375} - \frac{1360516D}{1875} + \frac{103522}{693(3D - 14)} - \frac{1720}{27(D - 1)} - \frac{9092}{15(D - 2)} \\
A_{6.2} &= +\frac{1489D^2}{25} - \frac{334171D}{375} + \frac{640}{3D} - \frac{8816}{21(2D - 5)} + \frac{239}{21(2D - 7)} \\

\end{align*}
\]
\[ +A_{6,3} \left( \frac{+24791D^2}{125} - \frac{-26566007D}{364768} - \frac{2390}{8} + \frac{63(2D - 7)}{53864} - \frac{-14271579}{22100(2D - 9)} + \frac{27(3D - 8)}{585166784} \right) \\
- \frac{225(D - 2)}{800736} - \frac{(D - 3)}{2043264} + \frac{45(D - 4)}{6208} - \frac{-3346875(5D - 14)}{2491573703} + \frac{365625(5D - 16)}{16210664} \\
+ \frac{3125(5D - 18)}{3125(5D - 22)} + \frac{15(D - 4)^2}{337500} \right) \\
+ B_{6,2} \left( \frac{+23D^2}{2} + \frac{-729D}{800736} + \frac{2261}{9(2D - 5)} + \frac{875}{320} + \frac{5056}{4} \right) \\
+ \frac{9(D - 4)}{9(D - 2)} - \frac{3(D - 2)^2}{3(D - 4)^2} + \frac{5401}{4} \right) \\
+ C_{6,1} \left( \frac{+23D^2}{2} + \frac{-729D}{800736} + \frac{2261}{9(2D - 5)} + \frac{875}{320} + \frac{5056}{4} \right) \\
+ \frac{9(D - 4)}{9(D - 2)} - \frac{3(D - 2)^2}{3(D - 4)^2} + \frac{5401}{4} \right) \\
+ A_{7,1} \left( \frac{+64D^2}{25} - \frac{-3084D}{6976} + \frac{2432}{4768} - \frac{362(2D - 7)}{40448} - \frac{1754}{290701952} - \frac{36841312}{268125(5D - 18)} \right) \\
- \frac{128}{128} + \frac{3(D - 2)^2}{3(D - 2)^2} + \frac{625}{625} \right) \\
- A_{7,2} \left( \frac{+83D^2}{20} - \frac{-1307D}{905568} - \frac{71585}{672(2D - 7)} + \frac{35360(2D - 9)}{35360(2D - 9)} - \frac{5(2D - 9)}{5(2D - 9)} \right) \\
+ \frac{955568}{4875(5D - 16)} + \frac{4875(5D - 16)}{125(5D - 18)} + \frac{51544}{2000} \right) \\
- A_{7,3} \left( \frac{+676D^2}{125} - \frac{-96D}{62816} + \frac{1824}{1888} - \frac{556}{5730624} + \frac{2925(2D - 9)}{30838784} + \frac{191648}{191648} \right) \\
+ \frac{225(D - 2)}{1875(5D - 22)} + \frac{15(D - 3)}{3(2D - 2)^2} + \frac{704}{3125} \right) \\
+ A_{7,4} \left( \frac{+313D^2}{250} - \frac{-10751D}{6724} + \frac{38}{750} + \frac{1807}{48387368} + \frac{159120(2D - 9)}{26422144} - \frac{228988}{28125(5D - 18)} \right) \\
- \frac{444664}{45(D - 2)} + \frac{128}{128} + \frac{5}{5654087} \right) \\
+ A_{7,5} \left( \frac{+313D^2}{250} - \frac{-10751D}{6724} + \frac{38}{750} + \frac{1807}{48387368} + \frac{159120(2D - 9)}{26422144} - \frac{228988}{28125(5D - 18)} \right) \\
- \frac{444664}{45(D - 2)} + \frac{128}{128} + \frac{5}{5654087} \right) \]
\[\begin{align*}
+ A_{8,1} & \left( + \frac{37071D - 1853}{5000} - \frac{4528}{45(D - 2)} - \frac{60}{7(D - 3)} - \frac{1474512}{21875(5D - 14)} + \frac{28125(5D - 16)}{125664} - \frac{5872}{13125(5D - 22)} + \frac{48(2D - 7)^2}{482997} \right) \\
+ B_{8,1} & \left( \frac{2D^3 - 25D^2 + 94D - 112}{(2D - 5)(2D - 7)(D - 2)^2} \right) \\
+ C_{8,1} & \left( \frac{2D^3 - 25D^2 + 94D - 112}{(2D - 5)(2D - 7)(D - 2)^2} \right)
\end{align*}\]

\[- A_{9,1} \left( + \frac{162D^2}{125} - \frac{576}{625} - \frac{128}{5(D - 2)} + \frac{30}{7(D - 3)} + \frac{197568}{3125(5D - 14)} - \frac{23232}{3125(5D - 16)} - \frac{51456}{3125(5D - 18)} + \frac{7008}{4375(5D - 22)} + \frac{350478}{3125} \right) \]

\[\begin{align*}
+ A_{9,2} & \frac{12(D - 4)(3D - 14)(15D^5 - 266D^4 + 1708D^3 - 5016D^2 + 6624D - 2944)}{5(D - 18)(5D - 14)(5D - 16)(5D - 16)(5D - 22)(D - 3)(D - 2)} \\
+ A_{9,4} & \left( + \frac{243D^2}{4000} - \frac{1539D}{2500} - \frac{973}{768(2D - 7)} - \frac{226533}{144400(2D - 9)} - \frac{16}{25(D - 2)} + \frac{9072}{10625(5D - 14)} - \frac{56144}{40625(5D - 16)} + \frac{14016}{3125(5D - 18)} + \frac{8936}{9375(5D - 22)} + \frac{108351}{80000} \right)
\end{align*}\]

\[X_{C\mathbb{C}_4}^g = \]

\[\begin{align*}
- B_{4,1} & \left( + \frac{752476D^2}{125} - \frac{123631802D}{1875} - \frac{93029552}{5525(2D - 9)} - \frac{1120}{9(3D - 10)} + \frac{385}{3(3D - 14)} - \frac{640}{9(3D - 8)} - \frac{93104}{25(D - 2)} - \frac{998}{3(D - 3)} + \frac{3093496}{45(D - 4)} - \frac{613534944}{3346875(5D - 14)} + \frac{1942431464}{31829588} + \frac{114439}{9375(5D - 18)} - \frac{28125(5D - 22)}{256} + \frac{256}{(D - 2)^2} \right) \\
- A_{5,1} & \left( + \frac{108328D^2}{125} - \frac{6789932D}{625} + \frac{1495117}{5525(2D - 9)} + \frac{224}{9(3D - 10)} - \frac{75}{(3D - 2)} - \frac{128}{3(3D - 3)} + \frac{135}{135(D - 4)} - \frac{286875}{288875(5D - 14)} + \frac{406250}{5(D - 16)} + \frac{3125}{3125(5D - 18)} - \frac{195901888}{45(4D - 2)} + \frac{378752}{15(4D - 4)} + \frac{34816}{5625} + \frac{195901888}{9375(5D - 22)} + \frac{14756354}{(4D - 4)^2} \right) \\
+ A_{5,2} & \left( + \frac{5268D^2}{375} - \frac{5132552D}{1875} - \frac{16310376}{5525(2D - 9)} - \frac{1376}{27(3D - 10)} + \frac{40625}{77(3D - 14)} - \frac{1600}{27(3D - 8)} - \frac{225}{225(D - 2)} + \frac{45}{45(D - 4)} + \frac{3346875}{3346875(5D - 14)} + \frac{1340625}{1340625(5D - 16)} - \frac{6431928}{3125(5D - 18)} - \frac{25131904}{5625(5D - 22)} + \frac{128}{(D - 2)^2} + \frac{34304}{15(4D - 4)^2} + \frac{1024}{(D - 4)^2} \right)
\]
\[
\begin{align*}
& -B_{5,1} \left( + \frac{28956D^2}{125} - \frac{1786684D}{625} - \frac{7136}{5(D - 2)} - \frac{44}{(D - 3)} + \frac{739904}{135(D - 4)} \right. \\
& + \frac{8126272}{84375(5D - 14)} + \frac{3125(5D - 16)}{3125(5D - 18)} - \frac{625(5D - 22)}{15(5D - 4)^2} + \frac{18944}{15(D - 4)^3} + \frac{25627244}{3125} \\
& + B_{5,2} \left( + \frac{46436D^2}{75} - \frac{2905928D}{375} - \frac{2912}{27(3D - 10)} + \frac{1600}{27(3D - 8)} - \frac{19792}{5(D - 2)} \\
& + \frac{221152}{27(D - 4)} + \frac{3375(5D - 14)}{5036} + \frac{625(5D - 16)}{532368} - \frac{625(5D - 18)}{125(5D - 22)} \\
& - \frac{53120}{9(5D - 4)(D - 4)^3} + \frac{25088}{675} \right) \\
& - A_{6,1} \left( + \frac{18928D^2}{375} - \frac{1201712D}{1875} + \frac{50924}{231(3D - 14)} - \frac{21832}{45(D - 2)} + \frac{10096}{45(D - 4)} \\
& - \frac{6678656}{196875(5D - 14)} + \frac{1159792}{61875(5D - 16)} - \frac{2932}{125(5D - 18)} - \frac{392}{9375(5D - 22)} \\
& + \frac{256}{2663584} \right) \\
& + A_{6,2} \left( + \frac{2026D^2}{25} - \frac{430946D}{375} - \frac{226533}{5525(2D - 9)} + \frac{616064}{225(D - 2)} + \frac{392}{3(D - 3)} \\
& + \frac{832}{15(D - 4)} + \frac{6648608}{6375(5D - 14)} - \frac{20564992}{73125(5D - 16)} + \frac{2663584}{5625(5D - 22)} - \frac{256}{(D - 2)^2} \\
& - \frac{10}{3(D - 3)^2} + \frac{8774231}{1875} \right) \\
& - A_{6,3} \left( + \frac{28662D^2}{125} - \frac{1856633D}{625} - \frac{4757193}{11050(2D - 9)} - \frac{242144}{75(D - 2)} + \frac{36}{(D - 3)} \\
& + \frac{17584}{15(D - 4)} + \frac{8105216}{159375(5D - 14)} + \frac{224944}{24375(5D - 16)} + \frac{265625}{625(5D - 18)} - \frac{3044096}{3125(5D - 22)} \\
& + \frac{2048}{256} + \frac{64132247}{625} \right) \\
& - A_{7,1} \frac{64(D - 3)(D^2 - 7D + 16)(5D^3 - 62D^2 + 236D - 288)}{(D - 2)(5D - 16)(5D - 14)(5D - 18)} \\
& + A_{7,2} \left( + \frac{49D^2}{10} - \frac{3181D}{50} - \frac{71585}{1008(2D - 7)} - \frac{226533}{17680(2D - 9)} + \frac{6224}{45(D - 2)} \\
& + \frac{862656}{14875(5D - 14)} - \frac{2925(5D - 16)}{26192} + \frac{26192}{125(5D - 18)} + \frac{50413}{200} \right) \\
& + A_{7,3} \left( + \frac{1676D^2}{125} - \frac{132544D}{625} - \frac{1312}{(D - 2)} + \frac{4504}{105(D - 3)} + \frac{4678464}{3125(5D - 14)} \\
& + \frac{241664}{241664} - \frac{66464}{9375(5D - 18)} - \frac{4375(5D - 22)}{14016} - \frac{28}{(D - 3)^2} + \frac{3149844}{3125} \right)
\end{align*}
\]
\[ -A_{7,4} \left( \frac{271D^2}{125} - \frac{48887D}{1875} + \frac{3197}{504(2D - 7)} - \frac{75511}{8840(2D - 9)} - \frac{2488}{45(D - 2)} \right) \\
+ \frac{1}{4} \left( \frac{74375(5D - 14)}{3(D - 3)} - \frac{778544}{365625(5D - 16)} - \frac{5256}{3125(5D - 18)} + \frac{196592}{28125(5D - 22)} \right) \\
- \frac{3}{3(D - 3)^2} + \frac{500}{47263} \right) \\
-A_{7,5} \left( \frac{271D^2}{125} - \frac{48887D}{1875} + \frac{3197}{504(2D - 7)} - \frac{75511}{8840(2D - 9)} - \frac{2488}{45(D - 2)} \right) \\
+ \frac{1}{4} \left( \frac{74375(5D - 14)}{3(D - 3)} - \frac{778544}{365625(5D - 16)} - \frac{5256}{3125(5D - 18)} + \frac{196592}{28125(5D - 22)} \right) \\
- \frac{3}{3(D - 3)^2} + \frac{500}{47263} \right) \\
+A_{9,1} \left( \frac{162D^2}{125} - \frac{13878D}{625} - \frac{576}{5(D - 2)} + \frac{30}{7(D - 3)} + \frac{197568}{3125(5D - 14)} \\
- \frac{3125(5D - 16)}{32232} - \frac{3125(5D - 18)}{51456} + \frac{4375(5D - 22)}{7008} + \frac{350478}{3125} \right) \\
-A_{9,4} \left( \frac{81D^2}{2000} - \frac{513D}{1250} - \frac{973}{707200(2D - 9)} - \frac{75511}{75(D - 2)} \\
+ \frac{112288}{10625(5D - 14)} - \frac{9344}{121875(5D - 16)} + \frac{3125(5D - 18)}{28125(5D - 22)} + \frac{36117}{40000} \right) \\
\]

\[ X_{C_{AN}}^g = \\
+ B_{4,1} \left( \frac{3872D}{25} + \frac{2860}{21(3D - 10)} - \frac{3387058}{81081(D - 1)} - \frac{220}{3(D - 2)} + \frac{74}{3(D - 3)} \right) \\
- \frac{81(D - 4)}{49984} + \frac{3375(5D - 14)}{48224} + \frac{1375(5D - 16)}{83328} + \frac{1197504}{1625(5D - 18)} + \frac{9217}{27(D - 1)^2} \\
+ \frac{16}{(D - 2)^2} + \frac{1}{(D - 3)^2} - \frac{6016}{9(D - 4)^2} - \frac{3584}{9(D - 4)^3} - \frac{187024}{375} \right) \\
-A_{5,2} \left( \frac{952D}{75} - \frac{272}{63(3D - 10)} - \frac{11224376}{81081(D - 1)} + \frac{320}{3(D - 2)} - \frac{1024}{81(D - 4)} \right) \\
+ \frac{85184}{10125(5D - 14)} - \frac{9984}{1375(5D - 16)} + \frac{93312}{1625(5D - 18)} - \frac{27(D - 1)^2}{27(D - 1)^2} - \frac{32}{(D - 2)^2} \\
- \frac{1024}{27(D - 4)^2} - \frac{16384}{1125} \right) \\
+B_{5,2} \left( \frac{616D}{75} - \frac{128}{63(3D - 10)} - \frac{1113304}{81081(D - 1)} - \frac{512}{81(D - 4)} + \frac{8832}{10125(5D - 14)} \right) \\
- \frac{1375(5D - 16)}{1625(5D - 18)} - \frac{46656}{2008} - \frac{512}{27(D - 1)^2} - \frac{41632}{27(D - 4)^2} - \frac{1125}{} \right) \\
+ A_{6,1} \left( \frac{4(5D - 16)(D - 4)}{3(D - 1)^2(D - 2)^2} \right) \]
\[-A_{7,1} \frac{32(D - 3)(D - 4)(5D^3 - 62D^2 + 236D - 288)}{(D - 1)(5D - 16)(5D - 14)(5D - 18)} \]

\[X_{C,F,N_2}^9 = \]

\[-B_{4,1} \left( + \frac{9344D}{25} - \frac{5440}{21(3D - 10)} - \frac{256}{15(3D - 8)} + \frac{81051008}{135135(D - 1)} - \frac{17024}{9(D - 4)} \right) + \frac{166656}{1375(5D - 16)} + \frac{2395008}{1625(5D - 18)} - \frac{2304}{(D - 4)^2} \frac{1024}{(D - 4)^3} \]

\[+ A_{5,2} \left( + \frac{7904D}{75} + \frac{640}{63(3D - 10)} - \frac{128}{9(3D - 8)} + \frac{17515744}{81081(D - 1)} - \frac{4096}{27(D - 4)} \right) + \frac{170368}{10125(5D - 14)} - \frac{19968}{1375(5D - 16)} - \frac{1625(5D - 18)}{186624(D - 4)^2} - \frac{1024}{375} \]

\[-B_{5,2} \left( + \frac{544D}{25} + \frac{640}{63(3D - 10)} + \frac{128}{9(3D - 8)} + \frac{5925536}{81081(D - 1)} - \frac{2048}{27(D - 4)} \right) - \frac{136576}{10125(5D - 14)} - \frac{17664}{1375(5D - 16)} - \frac{1625(5D - 18)}{93312(D - 4)^2} - \frac{512}{512} \]

\[+ A_{7,1} \frac{64(D - 3)(D - 4)(5D^3 - 62D^2 + 236D - 288)}{(D - 1)(5D - 16)(5D - 14)(5D - 18)} \]

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