The immune system is remarkable in its ability to produce antibodies (Abs) with virtually any specificity from a limited repertoire of germ line precursors. Although the contribution of sequence diversity to this molecular recognition has been studied for decades, recent models suggest that protein dynamics may also broaden the range of targets recognized. To characterize the contribution of protein dynamics to immunological molecular recognition, we report the sequence, thermodynamic, and time-resolved spectroscopic characterization of a panel of eight Abs elicited to the chromophoric antigen 8-methoxypyrene-1,3,6-trisulfonate (MPTS). Based on the sequence data, three of the Abs arose from unique germ line Abs, whereas the remaining five comprise two sets of siblings that arose by somatic mutation of a common precursor. The thermodynamic data indicate that the Abs recognize MPTS via a variety of mechanisms. Although the spectroscopic data reveal small differences in protein dynamics, the anti-MPTS Abs generally show similar levels of flexibility and conformational heterogeneity, possibly representing the convergent evolution of the dynamics necessary for function. However, one Ab is significantly more rigid and conformationally homogeneous than the others, including a sibling Ab from which it differs by only five somatic mutations. This example of divergent evolution demonstrates that point mutations are capable of fixing significant differences in protein dynamics. The results provide unique insight into how high affinity Abs may be produced that bind virtually any target and possibly, from a more general perspective, how new high affinity Abs may be produced that bind virtually any specificity from a limited repertoire of germ line precursors. The results support emerging models for the evolution of novel protein function.

The ability to rigorously characterize the dynamics of a chromophore’s environment has been revolutionized by ultrafast nonlinear optical methods, such as three-pulse photon echo peak shift (3PEPS) spectroscopy (37–39). 3PEPS spectroscopy...
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![Structure of MPTS](image)

relies on the application of a series of three short laser pulses. The first two pulses inscribe a phase grating in the sample. After a variable time delay, referred to as the population time, a third pulse triggers rephasing of the sample, but only chromophores that have an unaltered phase can contribute to rephasing and thus emit an echo signal. Fluctuations of the environment randomize the phase of the chromophore, which, with increasing population time, results in fewer chromophores contributing to rephasing and an apparent peak shift of the echo signal, commonly described by \( M(t) \), the energy gap autocorrelation function. \( M(t) \) is conveniently represented in the frequency domain with the corresponding spectral density, \( \rho(\omega) \) (40), which represents the frequency and amplitude of the environment's motions. Additionally, dynamics that appear static on the time scale of the 3PEPS experiment are manifest as a non-zero asymptote of the peak shift, thus making it possible to detect and quantify static heterogeneity (39, 40).

We have developed a system that allows for the use of 3PEPS to characterize Ab dynamics (8–11, 41). To achieve this, we immunized mice with different chromophore-protein conjugates and isolated Abs that evolved to bind the chromophore. Although Abs generally evolve to bind protein targets, small molecule targets (or antigens (Ags)) have played a central role in the study of Ab structure and function, and the chromophoric Ags are expected to bind as would any other. The Ab-Ag complexes were characterized using 3PEPS and transient grating (TG) spectroscopy, and motions on several distinct sets of time scales were observed. For one anti-fluorescein (Fl) Ab, we also characterized protein dynamics as a function of somatic evolution (affinity maturation) (10, 11). We found that evolution significantly restricted the observed dynamics, consistent with the evolution of flexible germ line Abs that are polyclonal and recognize their Ags via an induced fit- or conformational selection-like mechanism, into more rigid Abs that more specifically recognize their Ags via a lock-and-key-like mechanism.

To further explore the range of dynamics inherent to an immune response and how the dynamics might be tailored by somatic evolution, we now report the detailed analysis of eight Abs that evolved in mice to bind the chromophore 8-methoxypyrene-1,3,6-trisulfonate (MPTS; Fig. 1). Sequence, thermodynamic, and dynamic characterization via 3PEPS and TG spectroscopy reveal the range in affinities and dynamics possessed by the antibodies that constitute the immune response to MPTS and also provide interesting examples of parallel, convergent, and divergent evolution.

EXPERIMENTAL PROCEDURES

Generation and Sequencing of Anti-MPTS Ab—Production of the panel of monoclonal antibodies was performed as described previously using an MPTS-keyhole limpet hemocyanin conjugate (8). Heavy-chain variable region (\( V_H \)) and light-chain variable region (\( V_L \)) gene sequences were also determined as described previously (41). In brief, cDNA was synthesized using reverse transcriptase (Ambion) from mRNA isolated from fresh hybridoma cells. A primer set specifically designed for mouse variable immunoglobulin regions (42, 43) was then used to PCR-amplify the \( V_H \) and \( V_L \) genes from cDNA. The resulting gene fragments were sequenced using standard procedures. Nucleotide sequences of the \( V_H \) and \( V_L \) genes are provided in supplemental Data Set S1.

Isothermal Titration Calorimetry—Isothermal titration calorimetric experiments were performed using an MCS-ITC or a VP-ITC calorimeter (GE Healthcare) and analyzed with Origin 6.0 (Microcal) software. Titrations were performed between 20 and 35 °C. In each titration, 19 consecutive injections of MPTS every 240 s followed a single, preliminary injection into an 8 \( \mu \)M solution of Ab (IgG) in phosphate-buffered saline (PBS), to a final concentration of 150 \( \mu \)M MPTS. A titration of MPTS into buffer alone was used to correct for the heat of dilution of the titrant. The resulting \( \Delta H^o \) and \( \Delta G^o \) values were used to calculate \( T \Delta S^o = \Delta H^o - \Delta G^o \) at each temperature and \( \Delta C_p \), from a linear fit to \( \Delta H^o(T) \). The total entropy was approximated as follows,

\[
\Delta S^o = \Delta S^o_{\text{conf}} - 8 \text{ cal K}^{-1} \text{mol}^{-1} + \Delta C_p \ln(T/385 \text{ K})
\]  

(Eq. 1)

where \( \Delta S^o_{\text{conf}} \) is the entropy change from changes of internal motions of protein and ligand (41).

3PEPS and Transient Grating Experiments—Samples containing 700 \( \mu \)M Ab (IgG) and 700 \( \mu \)M MPTS in PBS were washed extensively with PBS by repeated dilution and microfiltration (10,000 molecular weight cut-off, Amicon Ultra, Millipore) to remove unbound MPTS, until the filtrate was free of MPTS as measured by UV-visible absorption. The experimental setup for 3PEPS and TG has been described previously (8). Briefly, the damped output of a Ti:sapphire regenerative amplifier system (Spectra Physics Spitfire, 5 kHz, 200-nJ pulses at 825 nm) was doubled in a 0.2-mm-thick type I BBO crystal and sent through a prism compressor, producing 416-nm pulses with ~45-fs duration, as determined by two-photon absorption in DMSO. Typical pulse energies at the sample were ~20 nJ/pulse. Beam splitters were then used to generate three beams of approximately equal intensity, two of which were then variably delayed using a pair of delay stages (PM6000, Kensington Laboratories). After passing through the delay stages, the three beams were arranged in an equilateral triangle of ~7 mm/side and focused with a 200-mm focal length plano-convex fused silica lens into a home-built spinning cell with a path length of 200 \( \mu \)m containing sample. The entire setup was maintained at 22 ± 1 °C. Proper masks were used to block the excitation light and allow only signal in the two phase-matching directions \( k_1 - k_2 + k_3 \) and \( -k_1 + k_2 + k_3 \) to pass to two silica photodiodes (DET110, Thor-
spectroscopic data analysis—tg data were fit to double exponential decays with origin 6.0 (microcal). parameters of the spectral densities were obtained from a simultaneous fit of the experimental 3peps data and the steady-state absorption spectra using mukamel’s response function formalism (40), as described in detail previously (41). in brief, the total spectral density, \( \rho(\omega) \), was calculated from the sum of the vibrations of the chromophore, \( \rho_{\text{MPTS}}(\omega) \), and the vibrations of the protein, \( \rho_{\text{Ab}}(\omega) \).

Intramolecular vibrational frequencies and excitation-induced displacements of mpts were obtained from quantum chemical calculations and validated by experimentally determined raman frequencies, as described previously (8). \( \rho_{\text{Ab}}(\omega) \) was modeled as the sum of an undamped brownian oscillator term (which we refer to as the brownian oscillator term, \( \rho_{\text{BO}}(\omega) \)) for the fastest motions and one or two overdamped brownian oscillator terms (which we refer to as the kubo terms, \( \rho_{\text{Ki}}(\omega) \)) for slower motions. specifically, the undamped brownian oscillator term,

\[
\rho_{\text{BO}}(\omega) = \frac{2 \lambda_{\text{BO}}}{\pi \omega} \left( \frac{\omega_{\text{BO}}^2 + \omega^2}{1 + \omega^2/\omega_{\text{BO}}^2} \right) (\text{eq. 2})
\]

was used to represent the inertial subpicosecond protein dynamics, where \( \lambda_{\text{BO}} \) is the reorganization energy, \( \omega_{\text{BO}} \) is the frequency, and \( \Gamma_{\text{BO}} \) is the damping constant of the brownian oscillator, and the kubo terms,

\[
\rho_{\text{Ki}} = \frac{\lambda_{\text{Ki}}}{\pi \omega} \frac{\tau_{\text{Ki}}}{1 + \omega^2/\tau_{\text{Ki}}^2} (\text{eq. 3})
\]

were used to represent the picosecond dynamics, where \( \lambda_{\text{Ki}} \) and \( \tau_{\text{Ki}} \) are the reorganization energy and time constant, respectively. signals for the 3peps experiment and the steady-state absorption spectra were calculated from the line-broadening function \( g(t) \) by using standard procedures (37, 38, 40). \( g(t) \) was calculated from \( \rho(\omega) \) using the following expression.

\[
g(t) = \int_0^\infty d\omega \rho(\omega) \sin(\omega t)
\]

\[
+ \int_0^\infty d\omega \rho(\omega) \coth\left( \frac{\hbar \omega}{2k_B T} \right) \left[ 1 - \cos(\omega t) \right] + \frac{(\Delta_{\text{inh}} t)^2}{2} (\text{eq. 4})
\]

The parameters in \( \rho_{\text{Ab}}(\omega) \) and the amount of static inhomogeneity, \( \Delta_{\text{inh}} \), in \( g(t) \) were varied using a least-squares simplex fit algorithm to obtain the best fit of both the experimental 3peps data (for population times from 10 fs to 100 ps) and the steady-state absorption spectra using a custom suite of c programs developed in our laboratory based on code kindly provided by dr. delmar Larsen (University of california, davis, CA). as previously observed (11), the amplitude (\( \lambda_{\text{BO}} \)), frequency (\( \omega_{\text{BO}} \)), and damping constant (\( \Gamma_{\text{BO}} \)) of the brownian oscillator term could not be fit unambiguously; thus, the values of \( \omega_{\text{BO}} \) and \( \Gamma_{\text{BO}} \) were fixed to average values from free fits to the 3peps decays. Simulation of the data with different values for these fixed terms resulted in only small variations in the remaining terms. all computations were performed on the 64-bit linux cluster garibaldi at the scripps research institute. a representative fit is shown in fig. 2.

results

production and sequence characterization of anti-mpts abs—anti-mpts monoclonal IgG abs were purified from ascites supernatants produced by standard techniques from swiss webster mice challenged with mpts conjugated to key-hole limpet hemocyanin (44). nucleotide sequences (supplemental material) were determined for each Ab from PCR-amplified V\textsubscript{H} and V\textsubscript{L} gene fragments using cDNA obtained from the corresponding hybridoma cell lines (42, 43). from the sequence data, we found that the eight mature anti-MPTS Abs arose from five different germ line precursor Abs (Figs. 3 and 4). Abs 4B2, 8H9, and 9D5 each arose from distinct germ line precursors. However, based on the nucleotide sequence of the stochastically generated V-D-J junction, it is evident that 3E6 and 7D5 are sibling Abs that diverged from a common germ line Ab and are thus differentiated only by the somatic mutations acquired after divergence. Similarly, Abs, 2E8, 3D3, and 10B8 are also siblings that diverged from a common germ line precursor.

thermodynamics of mpts binding—the MPTS dissociation constant (\( K_{\text{d}} \)) and binding enthalpy (\( \Delta H^\circ \)) were determined for
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Spectroscopic Characterization of MPTS-Ab Complexes—
The steady-state absorption spectra show that each Ab binds MPTS in a unique environment that is distinct from free solvent (Fig. 5 and Table 2). Relative to MPTS free in solution (PBS, pH 7.4), Ab binding causes a red shift in the absorption maximum of 3–6 nm and, with the exception of Ab 4B2, a reduction of the absorption line width. Ab 9D5, which has the most positive $\Delta S^0$, also shows the narrowest line width. Interestingly, MPTS bound to the sibling Abs 3E6 or 7D5 exhibited distinctly different absorption maxima of 406.6 and 408.8 nm, respectively, suggesting that they evolved to bind MPTS rather differently. In contrast, despite the differences in binding thermodynamics, MPTS bound to any of the triplet sibling Abs, 2E8, 3D3, or 10B8, exhibited nearly identical absorption maxima and line widths.

TG experiments were used to determine the excited state lifetime of MPTS when bound to the anti-MPTS Abs (Table 3). For all Abs, at least two exponential functions were required to fit the TG decay: a shorter time constant of 0.13–3.1 ps and a longer time constant of 84–287 ps. The decay component with the longer time constant is attributed to the excited state lifetime of MPTS when bound to the anti-MPTS Abs (Table 3).

Each Ab by isothermal titration calorimetry and were used to calculate the binding-induced changes in free energy ($\Delta G^0$), entropy ($\Delta S^0$), and heat capacity ($\Delta C_P$) (Table 1). Generally, $K_D$ and $\Delta G^0$ values ranged from 20 to 450 nM and from −8.7 to −10.5 kcal/mol, respectively. Large variations were also observed among the constituent $\Delta H^0$ and $\Delta T S^0$ values, which varied by 4.4 and 5.6 kcal/mol, respectively. Interestingly, Ab 10B8 binds MPTS with 5-fold and 2-fold lower affinity than its siblings 2E8 and 3D3, respectively, due to differences in $\Delta S^0$ values, which along with the absorption spectra discussed above, further suggests that MPTS experiences a different microenvironment in the two Abs.
3PEPS decays were measured up to 400 ps population time, and $M(t)$ and $\rho_{\text{Ab}}(\omega)$ were calculated from the data (see “Experimental Procedures”). As previously observed for the Ab-Fl complexes (41), we found 3PEPS decays for the Ab-MPTS complexes on three time scales (Fig. 6 and Table 4): a sub-100 fs decay corresponding to inertial protein or water motions; a decay component on a 3–10 ps time scale that is assigned to subconformational dynamics (11, 41); and a longtime signal offset (i.e. corresponding to a decay on time scales longer than the experimental time window) that is assigned to conformational dynamics (11, 41). In addition, sibling Abs 3E6 and 7D5 as well as Ab 8H9 show an additional decay component with time constants of ~1.2 and 0.5 ps, respectively. This additional decay could result from a second tier of subconformational protein dynamics or from motions of proximal water molecules whose motions are somewhat restricted relative to bulk water (a 500 fs decay component was observed previously for free MPTS in water (8)). However, as discussed above, the same three Abs also showed an atypical ~3 ps component in the TG decay. It is thus more likely that the additional 3PEPS decay component for Abs 3E6, 7D5, and 8H9 results from spectral shifting due to excited state population dynamics. The 3PEPS decay of Ab 4B2, which shows an unusually broad absorption spectrum and an unusually low $\Delta G^{\circ}$, both of which

### TABLE 1

Thermodynamic parameters at $T = 298$ K

| $K_p$ ($\text{mM}^{-1}$) | $\Delta G^{\circ}$ (kcal/mol) | $\Delta H^\circ$ (kcal/mol) | $T\Delta S^{\circ}$ (kcal/mol) | $\Delta C_p^{\circ}$ (kcal/mol) | $T\Delta S_{\text{conf}}^{\circ}$ (kcal/mol/K) |
|---------------------|-----------------------------|-----------------------------|-----------------------------|-----------------------------|-----------------------------|
| 4B2                 | 65 ± 48                     | $-9.9 \pm 0.4$              | $-11.32 \pm 0.20$           | $-1.5 \pm 0.2$              | $-0.158 \pm 0.033$          | $-9.9 \pm 2.8$              |
| 8H9                 | 276 ± 65                    | $-9.0 \pm 0.1$              | $-12.08 \pm 0.07$           | $-3.1 \pm 0.1$              | $-0.126 \pm 0.012$          | $-10.0 \pm 1.0$             |
| 9D5                 | 21 ± 6                      | $-10.5 \pm 0.2$             | $-9.65 \pm 0.07$            | $0.9 \pm 0.1$               | $-0.128 \pm 0.011$          | $-6.1 \pm 1.0$              |
| 2E8                 | 46 ± 10                     | $-10.1 \pm 0.1$             | $-12.94 \pm 0.02$           | $-2.9 \pm 0.1$              | $-0.145 \pm 0.003$          | $-10.5 \pm 0.3$             |
| 3D3                 | 95 ± 10                     | $-9.6 \pm 0.1$              | $-10.91 \pm 0.04$           | $-1.3 \pm 0.1$              | $-0.155 \pm 0.006$          | $-9.9 \pm 0.5$              |
| 10B8                | 209 ± 15                    | $-9.2 \pm 0.1$              | $-11.77 \pm 0.03$           | $-2.6 \pm 0.1$              | $-0.128 \pm 0.005$          | $-11.2 \pm 0.4$             |
| 3E6                 | 168 ± 35                    | $-9.3 \pm 0.1$              | $-14.04 \pm 0.04$           | $-4.7 \pm 0.1$              | $-0.112 \pm 0.007$          | $-10.0 \pm 0.6$             |
| 7D5                 | 450 ± 110                   | $-8.7 \pm 0.1$              | $-14.09 \pm 0.07$           | $-5.4 \pm 0.1$              | $-0.105 \pm 0.012$          | $-10.4 \pm 1.0$             |

### TABLE 2

Absorption maxima and line widths of MPTS in water or bound to an anti-MPTS Ab

| $\lambda_{\text{max}}$ (nm) | fwhm (nm) |
|----------------------------|-----------|
| Water                      | 403.7     | 20.0    |
| 4B2                        | 407.1     | 22.3    |
| 8H9                        | 409.5     | 18.5    |
| 9D5                        | 408.5     | 17.8    |
| 2E8                        | 408.6     | 18.2    |
| 3D3                        | 408.9     | 18.0    |
| 10B8                       | 408.9     | 18.0    |
| 3E6                        | 406.6     | 19.1    |
| 7D5                        | 408.8     | 18.8    |

### TABLE 3

Parameters for multiexponential fits to transient grating data

| $a_1$ | $\tau_1$ | $a_2$ | $\tau_2$ |
|-------|----------|-------|----------|
| %     | ps       | %     | ps       |
| 4B2   | $29 \pm 12$ | $130 \pm 30$ | $100$ | $287 \pm 37$ |
| 8H9   | $42.0 \pm 1.2$ | $3150 \pm 130$ | $58.0 \pm 1.9$ | $234 \pm 16$ |
| 9D5   | $54.1 \pm 8.8$ | $224 \pm 15$ | $100$ | $111 \pm 2$ |
| 2E8   | $22.2 \pm 0.9$ | $440 \pm 60$ | $77.8 \pm 0.1$ | $136 \pm 9$ |
| 3D3   | $15.1 \pm 1.2$ | $650 \pm 250$ | $84.9 \pm 2.2$ | $84 \pm 9$ |
| 10B8  | $15.5 \pm 0.1$ | $410 \pm 80$ | $84.5 \pm 0.3$ | $96 \pm 1$ |
| 3E6   | $40.3 \pm 0.8$ | $2910 \pm 50$ | $59.7 \pm 1.9$ | $89 \pm 1$ |
| 7D5   | $35.3 \pm 1.1$ | $2840 \pm 170$ | $64.7 \pm 2.3$ | $228 \pm 6$ |

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**FIGURE 5.** Absorption spectra of Ab-MPTS complexes.

**FIGURE 6.** 3PEPS decays for the Ab-MPTS complexes. Lines are fits to the data.
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### TABLE 4

Fit parameters for 3PEPS data

|    | $\lambda_{R0}$ | $\omega_{R0}^a$ | $\Gamma_{R0}^a$ | $\lambda_{K1}$ | $\tau_{K1}$ | $\lambda_{K2}$ | $\tau_{K2}$ | $\Delta_{\text{inh}}$ |
|----|-----------------|-----------------|-----------------|----------------|-------------|----------------|-------------|----------------|
| 4B2 | 699 ± 60        | 550             | 320             | 73 ± 19        | 448 ± 163   | 53 ± 23        | 4.1 ± 1.4   | 172 ± 34       |
| 8H9 | 469 ± 19        | 550             | 320             | 49 ± 12        | 68 ± 9      | 6 ± 2          | 153 ± 7     |
| 9D5 | 560 ± 13        | 550             | 320             | 49 ± 12        | 68 ± 9      | 6 ± 2          | 153 ± 7     |
| 2E8 | 527 ± 20        | 550             | 320             | 61 ± 5         | 4.8 ± 0.7   | 416 ± 5        | 116 ± 5     |
| 3D3 | 489 ± 14        | 550             | 320             | 70 ± 6         | 6.2 ± 0.5   | 141 ± 13       |             |
| 10B8| 492 ± 6         | 550             | 320             | 84 ± 10        | 1208 ± 102  | 24 ± 6         | 4.6 ± 1.5   | 141 ± 7        |

* Parameter fixed during fit.

### TABLE 5

Normalized reorganization energies

|       | $\lambda_{\text{inh}}$ | $\lambda_{K1}$ | $\lambda_{K2}$ | $\Delta_{\text{inh}}^a$ |
|-------|------------------------|----------------|-----------------|-------------------------|
| 4B2   | 83.4 ± 7.1             | 8.1 ± 1.5      | 85.0 ± 0.4      |                         |
| 8H9   | 73.0 ± 2.9             | 8.2 ± 3.6      | 7.4 ± 0.1       |                         |
| 9D5   | 84.1 ± 2.0             | 7.4 ± 1.9      | 8.5 ± 0.1       |                         |
| 2E8   | 85.1 ± 3.3             | 10.9 ± 1.5     | 4.0 ± 0.2       |                         |
| 3D3   | 83.9 ± 2.4             | 10.5 ± 0.8     | 5.6 ± 0.1       |                         |
| 10B8  | 80.6 ± 1.1             | 11.5 ± 0.9     | 7.9 ± 0.2       |                         |
| 3E6   | 71.8 ± 4.3             | 8.9 ± 2.3      | 8.4 ± 0.3       |                         |
| 7D5   | 79.8 ± 1.1             | 9.1 ± 1.5      | 7.4 ± 0.1       |                         |

* $\lambda_{\text{inh}} = \Delta_{\text{inh}}/2k_B T$, $T = 298$ K.

### FIGURE 7

Spectral densities of Ab-MPTS complexes.

The immune system has long served as a paradigm for the study of molecular recognition and its evolution. In this case, protein dynamics may play an especially key role because the number of foreign molecules that must be recognized far exceeds the potential number of different Abs, which is limited by available B-cells (32). We previously characterized a panel of anti-Fc Abs (41) as well as the dynamics of one as a function of its somatic evolution, which revealed that evolution significantly rigidified the Ab and restricted its conformational heterogeneity (10, 11). To further characterize the range of dynamics possible within the immunoglobulin fold common to all Abs and how it might be tailored by somatic evolution, we raised a panel of eight Abs to the chromophoric Ag MPTS. Sequencing revealed that a wide variety of germ line Abs may serve as suitable starting points for the evolution of high affinity MPTS binders. Generally, the immunoglobulin fold of an Ab is composed of two chains, $\text{V}_L$ and $\text{V}_{H^*}$, with the combining site composed of three hypervariable loops, or complementarity-determining regions (CDRs), contributed from each chain. As is common with Abs, only $\text{V}_L$ CDR1 and $\text{V}_{H^*}$ CDR3 show variations in length among the anti-MPTS Abs, with the former being 10–16 amino acids in length and the latter being 7–13 amino acids in length. In addition to being the longest, the 4B2 $\text{V}_{H^*}$ CDR3 also contains the largest number of charged residues (one Arg, one Glu, and three Asp residues), whereas the $\text{V}_L$ CDR3 loops of 8H9, 7D5, and 3E6 are the shortest and, in the case of the 7D5 and 3E6, also the most hydrophobic. Interestingly, we also found two groups of sibling Abs, Abs that evolved via somatic mutation from a common germ line precursor (Fig. 4). Thus, this panel of Abs provides an opportunity to study the diversity of protein dynamics available within the immune response to a given Ag and potentially how it is tailored by somatic evolution for molecular recognition.

**Antibody Dynamics and Molecular Recognition**—The anti-MPTS Ab evolved to bind MPTS with a relatively broad range in affinities, with $K_{D^*}$ values ranging from 20 nM (Ab 9D5) to 450 nM (Ab 7D5). These variations generally result from relatively larger variations in $\Delta H^O$ and $\Delta S^O$, which somewhat offset each other (Fig. 8). Such enthalpy-entropy compensation has been observed previously with other Abs (51), as well as with other ligand-binding proteins (52). Although the entropy differences should be interpreted cautiously because the calculated $\Delta S^O_{\text{cont}}$ values suggest that, with the exception of Ab 9D5, the differences in binding-induced changes in conformational entropy are smaller, the data clearly demonstrate that a threshold affinity is achieved via different mechanisms of molecular recognition.
As an Ab-MPTS complex fluctuates, the changing environment induces shifts in the MPTS transition frequency. In a 3PEPS experiment, this leads to dephasing and a decay in the 3PEPS signal. The time constants of the decay components observed in $M(t)$ are related to the frequency of motions within or between discrete states of the protein, whereas the corresponding amplitudes report on the amount of phase space sampled by a particular decay component. As was observed with anti-Fl Abs (41), the distinct time scales of motion observed in the Ab-MPTS complexes are not consistent with a continuum of motions within the Ab, but rather are reminiscent of Frauenfelder’s model of a hierarchical energy landscape wherein a protein exists in a limited number of conformations, each consisting of a large number of conformational substates (53, 54). Thus, we interpret the distinct time scales of motion observed in the Ab-MPTS complexes as different ranks of protein fluctuations corresponding to fluctuations between conformations as well as between and within conformational substates.

The rigidity of any material may be considered in terms of elastic and inelastic contributions. Rigid materials possess a higher level of elastic relative to inelastic motions. When described with a hierarchical energy landscape, a more rigid protein thus possesses higher levels of elastic fluctuations within conformational substates and, due to relatively higher energy barriers, lower levels of inelastic fluctuations between conformational states and substates (Fig. 9). In a 3PEPS experiment, the fast, inertial dynamics (i.e., motions within a conformational substate) are described by $\lambda_{BO}$ (amplitude), $\omega_{BO}$ (frequency), and $\Gamma_{BO}$ (damping). Little variation in elasticity, as measured with $\lambda_{BO}$, was observed among the anti-MPTS Abs, and it is probably determined largely by the motions of the charged side chains that are invariantly present within any protein.

The inelasticity of the Ab-MPTS complexes results from diffusive motion across barriers between different substates that are accessible on the time scale of the 3PEPS experiment, described by $\lambda_K$ (amplitude) and $\tau_K$ (time scale), as well as from diffusive motion between distinct conformational states that is slow on the 3PEPS time scale and thus manifest as static (conformational) inhomogeneity, $\lambda_{inh}$. The 3PEPS decays of the 3E6, 7D5, and 8H9 complexes show a component that occurs on the 0.5–1 ps time scale, which probably reflects excited state population dynamics because it also appears to be present in the TG signals. However, each Ab also shows a slower decay in its 3PEPS signals that is not present in its TG decay or with MPTS free in solution, and they are thus assigned to protein motions. The time constants for these diffusive motions ($\tau_{K2}$) generally ranged between 3 and 6 ps. The exception is Ab 3E6, for which significantly slower diffusive motion was observed ($\tau_{K2} = 13$ ps). Previous studies of dyes in small molecule solvents showed diffusive motions on the 1 ps time scale (38, 55), a previous study of eosin bound to lysozyme identified diffusive motions on the 7 ps time scale (56), and five of six of the previously reported Ab-Fl complexes showed diffusive motions on the 0.6–2.9 ps time scale, with one showing a 5.2 ps time scale diffusive motion (41). Using the Eyring model (57), the observed time constants translate into barrier heights of 1.8–2.2 kcal/mol for each Ab other than 3E6 and into 2.6 kcal/mol for Ab 3E6.

The static inhomogeneity observed with all but the 3E6-MPTS complex ranged between 101 and 172 cm$^{-1}$. By comparison, $\Delta_{inh}$ values for small molecule solvent are zero, whereas values of $\sim 500$ cm$^{-1}$ have been observed from the dye IR144 in acrylic glasses, which are expected to possess a larger range of static environments (58, 59), and the panel of six previously characterized Ab-Fl complexes showed $\Delta_{inh}$ values that ranged from 136 to 301 cm$^{-1}$. Ab 3E6 was again the exception because the $\Delta_{inh}$ value associated with this complex was within experimental error of zero. Along with the slower diffusive motions, this suggests that the 3E6-MPTS complex populates fewer conformations than the other Abs and that those conformations are separated by larger barriers.

**Antibody Evolution**—The majority of the anti-MPTS Abs evolved from different germ line precursors, and although differences in molecular recognition are apparent from the thermodynamic data, the flexibility and conformational heterogeneity of the Abs are generally remarkably similar. Moreover, the dynamics of the anti-MPTS Abs are similar to those of the previously characterized anti-Fl Abs (41). Although the germ lines of the anti-MPTS Abs have not yet been characterized, based on previous results (10, 11), it is likely that they are more distinct. If this is indeed the case, then the common dynamics may be the result of convergent evolution. Indeed, because germ line Abs are often broadly polyspecific and affinity-matured Abs are exquisitely specific, and because specificity is logically associ-
Antibody Dynamics

Ate with rigidity and conformational homogeneity, it seems likely that one role of somatic mutation is to evolve rigidity. Such mutations might act, for example, by introducing hydrogen bonds or packing interactions that cross-link the CDR loops or the strands of the β-sheets that anchor the loops, as has already been suggested for an anti-Fl Ab (11). Whatever the mechanism, selecting for rigidity, as a consequence of its effect on specificity, might be a natural consequence of the competition between B cells for limiting Ag during the evolution process in an environment full of other potential binding partners (as any biological compartment is certain to be).

Abs 2E8, 3D3, and 10B8 constitute a triplet set of siblings that evolved from a common germ line Ab. By maximum parsimony, comparison of their sequences allows for the identification of the somatic mutations. The sequences reveal that Ab 3D3 retains its germ line V_L sequence, whereas it acquired two V_H somatic mutations, Q39L and G56D. In contrast, Ab 2E8 acquired a single V_L mutation, S31N, during affinity maturation and two V_H mutations, S31N and W33L. Likewise, Ab 10B8 acquired a single V_L mutation, M33I, and two V_H mutations, S31N and W33L. Although the majority of these somatic mutations (six of eight) are in CDR loops, and some differences in affinity for MPTS are again apparent among the siblings, their flexibility and conformational heterogeneity are remarkably similar. In this case, the common dynamics are the result of parallel evolution.

Abs 3E6 and 7D5 constitute a set of twin siblings that also evolved from a common germ line Ab. Because only two sequences are available, it is not possible to identify the germ line sequence, but it was possible to identify five sites where mutations were introduced (i.e. the sites that differ). The V_L chains of the Abs differ at position 17 (Lys in 7D5; Glu in 3E6), position 78, (Met in 7D5; Leu in 3E6), and position 95 (Pro in 7D5; Ser in 3E6). The V_H chains differ at positions 35 (Asn in 7D5; Ser in 3E6) and 56 (Glu in 7D5; Asp in 3E6). Three of these sites are within CDR loops (Pro/Ser-95, Asn/Ser-35, and Glu/Asp-56). Abs 3E6 and 7D5 show similar values of both ΔS° and ΔC°sp which suggests that complex formation involves less desolvation than with the other Abs. Nonetheless, Abs 3E6 and 7D5 show differences in MPTS affinity and, unlike the triplet siblings discussed above, significant differences in dynamics; Ab 3E6 binds MPTS with higher affinity than 7D5 and it is uniquely rigid and conformationally homogeneous compared with all of the other Abs characterized. It is possible that the relative flexibility of Ab 7D5 in shared with its germ line Ab, and if so, then the rigidity of Ab 3E6 would be further evidence of selection of rigidity and conformational homogeneity during somatic evolution, as already demonstrated with an anti-Fl Ab (11). Although, as mentioned above, it is not possible to rigorously identify the germ line precursor of the twin siblings, comparisons with homologous germ line genes suggest that they probably have a similar number of somatic mutations. Thus, regardless of whether the increased rigidity was itself selected for, these sibling Abs provide a clear demonstration of the power of divergent evolution to tailor protein dynamics and thus provide an opportunity for altered molecular recognition.

CONCLUSION

Eight anti-MPTS monoclonal Abs were generated and characterized, and although they clearly recognize MPTS via different mechanisms, they all show a similar level of flexibility and conformational heterogeneity that is also similar to that of a panel of anti-Fl Abs (41). It is interesting to speculate that this may reflect the levels of affinity and specificity required for immune function. The data also demonstrate that somatic evolution may have a pronounced effect on protein dynamics, which further supports a connection between rigidity, specificity, and Ab evolution. Future structural efforts will focus on elucidating the mechanisms by which the somatic mutations rigidify the Ab-combining site.

Inherently flexible germ line Abs and point mutations that rigidify them may play an important role in the evolution of molecular recognition within the immune response by maximizing the likelihood that a germ line Ab with the required activity is present and then tailoring dynamics for specificity, so that they may be produced at the high levels required to eradicate an infection without causing autoimmunity. Some level of flexibility and associated polyspecificity is inherently present in all proteins, and point mutations that favor conformations with activities and selectivities that satisfy a given selection pressure probably play an important role in the evolution of other proteins, such as enzymes, transcription factors, receptors, and components of signaling pathways.

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