THE RADIAL AND AZIMUTHAL PROFILES OF Mg\textsc{ii} ABSORPTION AROUND 0.5 < z < 0.9 zCOSMOS GALAXIES OF DIFFERENT COLORS, MASSES, AND ENVIRONMENTS

R. Bordoloi\textsuperscript{1}, S. J. Lilly\textsuperscript{1}, C. Knobel\textsuperscript{1}, M. Bolzonella\textsuperscript{2}, P. Kampczyk\textsuperscript{1}, C. M. Carollo\textsuperscript{1}, A. Iovino\textsuperscript{3}, E. Zucca\textsuperscript{2}, T. CONTINI\textsuperscript{4,5}, J.-P. Kneib\textsuperscript{6}, O. LE FEVRE\textsuperscript{6}, V. MAINIERI\textsuperscript{6}, A. RENZINI\textsuperscript{6}, M. SCODEGGO\textsuperscript{6}, G. ZAMORANI\textsuperscript{6}, L. BALESTRA\textsuperscript{10}, S. BARDELLI\textsuperscript{2}, A. BONGIORNO\textsuperscript{10}, K. CAPUTI\textsuperscript{11}, O. CUCCIATTI\textsuperscript{12}, S. DE LA TORRE\textsuperscript{11}, L. DE RAEL\textsuperscript{11}, B. GARILO\textsuperscript{9,6}, K. KOVA\textsuperscript{9,13}, F. LAMAREILLE\textsuperscript{4,5}, J.-F. LE BORGNE\textsuperscript{4,5}, V. LE BRUN\textsuperscript{6}, C. MAIER\textsuperscript{5}, M. MIGNOLI\textsuperscript{2}, R. PELLO\textsuperscript{4,5}, Y. PENG\textsuperscript{1}, E. PEREZ MONTERO\textsuperscript{4,5,14}, V. PRESOTTO\textsuperscript{2}, C. SCARLATA\textsuperscript{13}, J. SILVERMAN\textsuperscript{16}, M. TANAKA\textsuperscript{16}, L. TASC\textsuperscript{6}, L. TRESSE\textsuperscript{6}, D. VERNANDI\textsuperscript{2}, L. BARNES\textsuperscript{1}, A. CAPPI\textsuperscript{2}, A. CIMATTI\textsuperscript{11}, G. COPPA\textsuperscript{10}, C. DIENER\textsuperscript{1}, P. FRANZETTI\textsuperscript{9}, A. KOEKEMOER\textsuperscript{18}, C. LÓPEZ-SANJUÁN\textsuperscript{6}, H. J. MCCrackEN\textsuperscript{19}

M. MORESCO\textsuperscript{17}, P. NAIR\textsuperscript{8}, P. OSCH\textsuperscript{1,20}, L. POZZETTI\textsuperscript{2}, AND N. WELIKALA\textsuperscript{21}

\textsuperscript{1} Institute for Astronomy, ETH Zürich, Wolfgang-Pauli-Strasse 27, 8093 Zürich, Switzerland; rongmonb@phys.ethz.ch
\textsuperscript{2} INAF Osservatorio Astronomico di Brera, Milan, Italy
\textsuperscript{3} INAF Osservatorio Astronomico di Bologna, Bologna, Italy
\textsuperscript{4} Institut de Recherche en Astrophysique et Planétologie, CNRS, 14, Avenue Edouard Belin, F-31400 Toulouse, France
\textsuperscript{5} IRAP, Université de Toulouse, UPS-OMP, Toulouse, France
\textsuperscript{6} Laboratoire d’Astrophysique de Marseille, CNRS/Aix-Marseille Université, 38 rue Frédéric Joliot-Curie, 13388 Marseille Cedex 13, France
\textsuperscript{7} European Southern Observatory, Garching, Germany
\textsuperscript{8} Dipartimento di Astronomia, Università di Padova, Padova, Italy
\textsuperscript{9} INAF-IASF Milano, Milano, Italy
\textsuperscript{10} Max Planck Institut für Extraterrestrische Physik, Garching, Germany
\textsuperscript{11} SUPA, The University of Edinburgh, Royal Observatory, Blackford Hill, Edinburgh EH9 1BD, UK
\textsuperscript{12} INAF-Osservatorio Astronomico di Trieste, Trieste, Italy
\textsuperscript{13} Max-Planck-Institut für Astrophysik, Garching, Germany
\textsuperscript{14} Instituto de Astrofísica de Andalucía, CSIC, Apdo correes 5004, 18080 Granada, Spain
\textsuperscript{15} School of Physics and Astronomy, University of Minnesota, Minneapolis, MN 55455, USA
\textsuperscript{16} Institute for the Physics and Mathematics of the Universe (IPMU), University of Tokyo, Kashiwanoha 5-1-5, Kashiwa-shi, Chiba 277-8568, Japan
\textsuperscript{17} Dipartimento di Astronomia, Università degli Studi di Bologna, Bologna, Italy
\textsuperscript{18} Space Telescope Science Institute, Baltimore, MD 21218, USA
\textsuperscript{19} Institut d’Astrophysique de Paris, UMR7095 CNRS, Université Pierre & Marie Curie, 75014 Paris, France
\textsuperscript{20} UCO/Lick Observatory, Department of Astronomy and Astrophysics, University of California, Santa Cruz, CA 95064, USA
\textsuperscript{21} Institut d’Astrophysique Spatiale, Batiment 121, CNRS & Université Paris Sud XI, 91405 Orsay Cedex, France

Received 2011 June 2; accepted 2011 August 24; published 2011 November 17

ABSTRACT

We map the radial and azimuthal distribution of Mg\textsc{ii} gas within ∼ 200 kpc (physical) of ∼ 4000 galaxies at redshifts 0.5 < z < 0.9 using co-added spectra of more than 5000 background galaxies at z > 1. We investigate the variation of Mg\textsc{ii} rest-frame equivalent width (EW) as a function of the radial impact parameter for different subsets of foreground galaxies selected in terms of their rest-frame colors and masses. Blue galaxies have a significantly higher average Mg\textsc{ii} EW at close galactocentric radii as compared to the red galaxies. Among the blue galaxies, there is a correlation between Mg\textsc{ii} EW and galactic stellar mass of the host galaxy. We also find that the distribution of Mg\textsc{ii} absorption around group galaxies is more extended than that for non-group galaxies, and that groups as a whole have more extended radial profiles than individual galaxies. Interestingly, these effects can be satisfactorily modeled by a simple superposition of the absorption profiles of individual member galaxies, assuming that these are the same as those of non-group galaxies, suggesting that the group environment may not significantly enhance or diminish the Mg\textsc{ii} absorption of individual galaxies. We show that there is a strong azimuthal dependence of the Mg\textsc{ii} absorption within 50 kpc of inclined disk-dominated galaxies, indicating the presence of a strongly bipolar outflow aligned along the disk rotation axis. There is no significant dependence of Mg\textsc{ii} absorption on the apparent inclination angle of disk-dominated galaxies.

Key words: galaxies: evolution – galaxies: groups: general – galaxies: high-redshift – intergalactic medium – ISM: jets and outflows – quasars: absorption lines

Online-only material: color figures

1. INTRODUCTION

Metal absorption lines, such as those of the Mg\textsc{ii} λλ 2796, 2803 doublet, in the spectra of background sources such as quasars, provide an important tracer of enriched gas that is otherwise very hard to detect. Mg\textsc{ii} absorption originates in photoionized gas at temperatures around T ∼ 10\textsuperscript{4} K (Bergeron & Stasińska 1986; Charlton et al. 2003) with neutral hydrogen column densities of N(H\textsc{i}) ≈ 10\textsuperscript{16}–10\textsuperscript{22} cm\textsuperscript{-2} (Churchill et al. 2000; Rigby et al. 2002; Rao et al. 2006). Optical spectrographs can detect the Mg\textsc{ii} doublet over a broad redshift range of z ∼ 0.3–2.5 and over the years there have been a large number of studies using QSO absorption lines to characterize the statistical properties of Mg\textsc{ii} absorbers.

These works have studied the distribution of column densities, the redshift evolution of number densities, and the kinematic signatures (see, e.g., Lanzetta et al. 1987; Sargent et al. 1988; Petitjean & Bergeron 1990; Steidel & Sargent 1992; Charlton...
It has been postulated that strong Mg\textsc{ii} absorption with normal, bright, field galaxies is by now well established (e.g., Churchill et al. 2005b and references herein). The Mg\textsc{ii} gas around galaxies is traced out to $\sim 100$ kpc with absorber covering fractions of 50\%–80\% (Chen et al. 2010a). The profile of the Mg\textsc{ii} halos around galaxies has also been studied quite extensively. Steidel (1995) showed that their sizes scale weakly with galaxy luminosity and inferred that these halos have quite sharp boundaries. Churchill et al. (2005a) searched for correlations between gas kinematics and galaxy orientation and found no significant correlation between the two for their small sample of absorbers. Kačprzak et al. (2008) suggested that covering fractions might be less that unity and that Mg\textsc{ii} halos around galaxies are patchy and might have non-symmetric geometric distribution.

However, despite such observational progress, the origin of the absorbing gas in the halos of galaxies is still widely debated. It has been postulated that strong Mg\textsc{ii} systems originate due to cool gas entrained in outflows from star-forming galaxies (Bouché et al. 2006; Bouché 2008; Ménard et al. 2009; Nestor et al. 2010), which is also supported by the blueshifts of Mg\textsc{ii} absorption in the spectra of individual galaxies (Weiner et al. 2009; Rubin et al. 2010). Others have suggested in falling gas as the source (Chen et al. 2010a; Chen & Tinker 2008; Kačprzak et al. 2010a), or a combination of both inflows and outflows (Chen et al. 2010b; Chelouche & Bowen 2010).

It has also been suggested that group galaxies may have different radial profiles in Mg\textsc{ii} absorption compared with isolated galaxies (Chen et al. 2010a). A number of mechanisms for this have been suggested, such as tidal tails and streams, etc. (Bowen et al. 1995; Churchill & Charlton 1999; Kačprzak et al. 2010a, 2010b). Recent observational studies have utilized close galaxy pairs and galaxy groups to study Mg\textsc{ii} absorption strength around such systems (Chen et al. 2010a; Nestor et al. 2007; Kačprzak et al. 2010a, 2010b) but these studies have generally been limited to small samples of absorbers. Most of the above investigations have utilized the spectra of bright quasars for which high-resolution, high signal-to-noise (S/N) spectra can be readily obtained. More recently, it has become practical to use the spectra of background galaxies that have been observed in large-scale redshift surveys (Steidel et al. 2010). The spectra of these are sufficiently noisy such that it is usually necessary to stack the spectra of many background galaxies at a given impact parameter $b$, so as to be able to detect the absorption signal. This approach differs from that of using quasar spectra in a number of significant ways.

1. The individual lines of sight probe regions that are of order 10 kpc across, or more, as compared with the 1 pc sight lines of quasar spectra.
2. Furthermore, the stacked spectra for a given range of $b$ yield a measure of the average absorption over an annular region around the foreground systems encompassing an area of several thousand kpc$^2$, i.e., about $10^7$ times larger in area as compared to a QSO sight line.
3. This average spectrum will integrate the contribution of individual absorption systems that have too low equivalent width (EW) to be detected in even high-S/N quasar spectra.
4. Some of the difficulties that are encountered in correctly associating foreground galaxies with individual absorption systems (e.g., due to the glare of the bright quasar or limitations in the depth of the associated galaxy follow-up) are mitigated by using the spectra of background galaxies around a well-defined set of foreground galaxies, rather than starting with a quasar spectrum and trying to identify nearby foreground galaxies responsible for the absorption.

This approach is inherently statistical, since generally the absorption around multiple foreground galaxies must be co-added. Against this, the number of spectra that are potentially available in large redshift surveys can be high enough to allow the statistical study of different sub-samples of foreground populations. This approach is orthogonal to most of the previous studies of Mg\textsc{ii} absorption systems, in the sense that those studies select the samples of galaxies from detected Mg\textsc{ii} absorption, comparing these with the general galaxy population. Our own approach takes a well-defined population of galaxies and then measures the absorbing gas around them. It should be noted that because we are summing (or averaging), the spectra across the spatial rather than the spectral domain, the derived EWs will be correct, independent of the degree of saturation of the absorption along any given line of sight. This is different from the case of absorption components which overlap spectrally along a given line of sight where saturation effects must be carefully considered. Of course, conversion of our mean EW to a mean column density along the annulus would need to consider the effects of saturation, but this will not be attempted in this paper.

The aim of this paper is to apply this approach to the available galaxy spectra in the zCOSMOS survey (Lilly et al. 2007). This is a large redshift survey that has been undertaken in the COSMOS field (Scoville et al. 2007). In particular we use the blue spectra of galaxies in zCOSMOS-deep, selected to have $z_{\text{phot}} > 1$ to probe the Mg\textsc{ii} absorption around foreground galaxies that have secure spectroscopic redshifts from zCOSMOS-bright in the $0.5 < z < 0.9$ range.

There are two major advantages of doing this study in the COSMOS field. First, the high density of spectroscopic redshifts yields good environmental information for the foreground galaxies in the form of a high fidelity group catalog (Knobel et al. 2009; C. Knobel et al. 2011, in preparation). Second, high-resolution Hubble Space Telescope (HST)/Advanced Camera for Surveys (ACS) images (Koekemoer et al. 2007) are available for all galaxies, enabling us to study the azimuthal dependence of Mg\textsc{ii} absorption relative to the disc axis for disc-dominated galaxies.

This paper is organized as follows. In Section 2 we first present the spectroscopic data sets that are used, and describe the selection criteria and the derivation of the final sample. In Section 3 we construct the radial profile of Mg\textsc{ii} absorption around foreground galaxies at $0.5 < z < 0.9$ and examine the dependence of this on the color, stellar mass, and group environment of the galaxies. In Section 3.1 we show directly that there is a clear mass dependence in the strength of Mg\textsc{ii} absorption around blue galaxies, and that, at a given mass, the absorption around blue galaxies is much stronger than around red ones. In Section 3.2, we show that the radial dependence of Mg\textsc{ii} absorption around group galaxies is considerably more extended than around isolated galaxies. We also construct the absorption profile around the group centers, and around the most massive members of the groups, and show that this is also more extended than around isolated galaxies. We show that these effects can however be fully explained by a simple model which superposes the normal absorption profiles of the individual members. In Section 3.3 we study the azimuthal dependence of Mg\textsc{ii} absorption around inclined disc galaxies.
and show that the absorption along the projected rotation axes of the disks is about three times as strong as that in perpendicular directions. Finally, in Section 3.4 we investigate the dependence of \( \Omega_{\text{m}} \) absorption on the apparent inclination angles of disk galaxies.

Throughout this paper, we use a concordance cosmology with \( \Omega_{\text{m}} = 0.25, \Omega_{\lambda} = 0.75, \) and \( H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1} \). Unless stated otherwise all magnitudes are given in the AB system. Because we are working with very low resolution spectra, all EWs are quoted integrating over both components of the Mg \( \text{II} 2796, 2803 \) doublet. They are therefore approximately twice as large as the EWs of the individual components.

2. SPECTROSCOPIC DATA

2.1. The zCOSMOS Redshift Survey

The zCOSMOS survey (Lilly et al. 2007) is a survey in the 2 deg\(^2\) COSMOS field (Scoville et al. 2007). The zCOSMOS survey is carried out with the VIMOS spectrograph on the ESO UT3 8 m Very Large Telescope. The survey itself is divided into two major components:

1. \text{zCOSMOS-bright}, which consists of spectra from approximately 20,000 flux limited \( I_{AB} < 22.5 \) galaxies over the full 2 deg\(^2\) COSMOS field. At this flux limit, the majority of the observed galaxies have redshifts in the range of \( 0 < z < 1.4 \). Observations in \text{zCOSMOS-bright} were obtained with the medium resolution grism using 1 arcsec slits, yielding a spectral resolution of \( R \sim 600 \) at 2.5 Å pixel\(^{-1}\). The average accuracy of individual redshifts has been demonstrated to be 110 km s\(^{-1}\) (Lilly et al. 2009). Stellar masses for these galaxies have been estimated by spectral energy distribution fitting, using the Hyperz mass code, a modified version of the \textit{photo-z} code \textit{Hyperz} (Bolzonella et al. 2000). We refer the reader to Bolzonella et al. (2010) for a detailed technical description of this mass estimation. The absolute magnitudes used in this study were computed as in Zucca et al. (2009).

2. \text{zCOSMOS-deep}, targets the central 1 deg\(^2\) of the field and consists of approximately 10,000 spectra of \( B < 25.25 \) color-selected galaxies which are expected to predominantly lie in the redshift range of \( 1.4 < z < 3 \) (Lilly et al. 2007; S. J. Lilly et al. 2011, in preparation). The spectra were obtained using the low resolution blue grism with a resolution \( R \sim 200 \) at 5.3 Å pixel\(^{-1}\). The spectra cover the wavelength range from 3500 Å to 7000 Å. For each of the objects observed spectroscopically, photometric redshifts are available from the excellent COSMOS photometry using a number of codes (e.g., Ilbert et al. 2009). The \textit{photo-z} accuracy for these objects for which reliable spectroscopic redshifts are available is \( \sigma_z \sim 0.0381(1+z) \). Further details about these two parts of zCOSMOS are given in Lilly et al. (2007, 2009) and S. J. Lilly et al. (2011, in preparation).

2.2. Foreground and Background Galaxy Selection

The first step in the analysis is to identify pairs of galaxies consisting of one foreground galaxy from \text{zCOSMOS-bright}, selected to lie at \( 0.5 < z_{\text{spec}} < 0.9 \), and one background galaxy from \text{zCOSMOS-deep} whose sight line passes within a specified projected radius (or impact parameter) \( b \) from the foreground galaxy. The foreground redshift range of \( 0.5 < z_{\text{spec}} < 0.9 \) is chosen so that the Mg \( \text{II} 2799 \) absorption at the foreground redshift lies within the spectral range of the \text{zCOSMOS-deep} spectra. For the foreground galaxies we only use galaxies with high-confidence redshift measurements, i.e., those with redshift Confidence Classes of 4.x, 3.x, 2.5, and 1.5. As shown in Lilly et al. (2009) these are 99% reliable. Based on the \textit{photo-z} estimates of the other galaxies, this secure sample is, in this redshift range, about 95% complete.

The background sample is selected from the \text{zCOSMOS-deep} survey. Because we do not need to know the exact redshift of the background object, provided only that we know it truly lies well behind the foreground one, we can use the spectra of background galaxies even if a confident redshift has not been secured. We therefore select the background galaxies using their \textit{photometric} redshifts to have \( z_{\text{phot}} > 1.0 \).

Figure 1 shows the redshift distributions of the foreground galaxies and the background galaxies, both of the full samples and of those galaxies that are used in this analysis because they form suitably close projected pairs. The total number of foreground galaxies between \( 0.5 < z < 0.9 \) is 7110. Of these, 3908 have at least one background galaxy projected within \( b \sim 200 \) kpc. These background galaxies comprise 5237 of the total potential background sample of 8529 galaxies. As expected, the samples that can be used for this study are a random set of the parent samples (Figure 1).

Clearly a major problem will occur if a “background” galaxy actually lies at the same redshift as the foreground galaxy, since then strong interstellar absorption in the former will be mistaken as absorption from the intergalactic medium around the latter. In the current study, this should be unlikely: most galaxies in the background sample lie at \( z_{\text{phot}} > 1.4 \), reflecting the use of the well-established \( gzK \) and \( ugr \) color selection criteria in \text{zCOSMOS-deep}. Random scattering to redshifts \( 0.5 < z < 0.9 \) is unlikely. Catastrophic failures in the \textit{photo-z} will generally involve degeneracies between \( z > 2 \) and \( z < 0.5 \). The loss of high-redshift galaxies with spuriously low photometric redshifts will be of no consequence except to reduce the size of the sample. The inclusion of low-redshift galaxies that are assigned a spuriously high \textit{photo-z} will merely dilute the signal from the foreground objects, since these sight lines do not in fact penetrate through the foreground system, but will not introduce a false signal. Using \text{zCOSMOS-deep} objects having secured
Figure 2. Foreground non-group and group galaxies divided into blue and red sub-samples as shown by the horizontal line computed from Equation (1). The left panel gives the selected non-group sample and the right panel gives the selected group sample. For both the panels the red points are for red galaxies and the blue points are for blue galaxies. The vertical lines on either panel are the boundaries chosen to divide the sample in mass as described in Section 3.2.

(A color version of this figure is available in the online journal.)

spectroscopic redshifts (i.e., zCOSMOS-deep galaxies with spectroscopic redshift Confidence Classes 4.x and 3.x), we estimate that the fraction of objects in this latter category to be about 3%. We have not removed these known interlopers, but have checked that none of them lie at the same redshift as the foreground galaxy. We associate to each foreground galaxy a stellar mass and corresponding absolute magnitude as described in the previous section. We divide our foreground galaxies into blue star-forming and red passive galaxies on the basis of their rest-frame \((u-B)\) color, which is a weak function of mass. The line dividing the blue and red galaxies is similar but not identical to that used in (Peng et al. 2010):

\[
(u - B)_{AB} = 0.98 + 0.075 \log \left( \frac{M}{10^{10} M_\odot} \right) - 0.18z, \tag{1}
\]

where \(M\) is the mass of the galaxy in question, for simplicity we use the average redshift \(z \sim 0.7\) of the sample as a whole. The color–mass division between blue and red galaxies is shown in Figure 2. The left-hand panel shows this division for non-group galaxies and the right-hand panel shows for the group galaxies.

For environmental information, we use the group catalog of C. Knobel et al. (2011, in preparation) to separate the foreground sample into group and non-group galaxies. Almost 25% of the initial sample of 7110 galaxies can be assigned to a group. We refer the reader to Knobel et al. (2009) for a detailed description of the group-finding algorithm and details of the zCOSMOS group catalog.

To study the azimuthal dependence on Mg \(\text{II}\) absorption line strength, we use the Zurich Estimator of Structural Types (ZEST) morphological classification (Scarlata et al. 2007) for each potential foreground galaxy. This is based on the \(HST/\text{ACS}\) F814W images of the COSMOS field. We first isolate those foreground galaxies that have a disk-dominated morphology, i.e., ZEST Type 2 excluding bulge-dominated systems, and further require that the galaxy has \(0 < b/a < 0.65\), corresponding to inclination angles of \(50 < i < 90\). This yields a sample of 595 foreground systems probed by 1134 background spectra within \(b < 200\) kpc.

2.3. Co-addition of Galaxy Spectra

The spectra of all background galaxies whose sight lines pass within a given region of a given set of foreground galaxies are then co-added as follows. The spectrum of each individual background galaxy is first shifted in wavelength to the rest frame of the foreground galaxy in question. This means that absorption at the redshift of the foreground galaxy will appear at the correct “rest wavelength,” i.e., 2799 Å in the shifted spectrum and will thereby allow co-addition of the spectra of background galaxies associated with foreground galaxies at different wavelengths. The continuum is then normalized by fitting the continuum with a running median filter of width 59 Å and then dividing the spectrum by it. To facilitate co-adding the spectra of different galaxies, this normalized spectrum is re-sampled onto a uniform grid in wavelength. The co-added spectra are produced by taking the median of the normalized re-sampled spectra. The median is chosen to reduce the sensitivity to absorption or emission features in the background galaxies, and to any other artifacts such as sky residuals, even though the effects of these should be random because of the range of redshifts of both foreground and background sources. We checked that other summation techniques, including straight averaging and average-sigma-clipping, produce very similar co-added spectra, but a marginally less clean continuum. The use of the median, plus the extended nature of the background source, means that the absorption in the co-added spectra represents a “typical” sight line rather than a true average including rare high-absorption sight lines. Three examples of the resulting co-added spectra are shown in Figure 3.

The spectral resolution of the zCOSMOS-deep spectra is very low \((R \sim 200)\), so the two components of the Mg \(\text{II}\) \(\lambda\lambda 2796, 2803\) doublet cannot be separated. The wavelength FWHM of almost 30 Å is much larger than the observed-frame 12 Å separation at the redshifts of interest. We therefore fit a single Gaussian profile around the Mg \(\text{II}\) rest-frame wavelength and measure the (rest-frame) absorption EW, integrating across both components of the doublet. In order to limit the number of free parameters in the fit, we determine the shape of this Gaussian (central wavelength and dispersion)
from a fit to the best absorption spectrum, which is obtained within 40 kpc of all foreground galaxies (shown in Figure 3, top spectra). This Gaussian profile is then used for all other spectra, with only the depth as a free parameter. We have checked that other integration schemes, e.g., simply summing the flux deficit through the region of the line, give indistinguishable results within the uncertainties.

The errors on the EW measurements are determined using a bootstrap approach. For each set of background spectra, a thousand co-added spectra are generated from random selections of the sample. The width of the distribution of measured EWs in these thousand spectra are taken as representative of the error in the original measurement, and should account for sample variance, continuum uncertainties, and profile fitting errors.

3. RESULTS

In the following sections we discuss the variation of Mg II absorption EW with different properties of the foreground galaxies, examining the radial profile, for the inclined disk galaxies, the azimuthal dependence and for disk galaxies, the dependence on apparent inclination.

3.1. Radial Dependence and Variation with Stellar Mass and Color

We first investigate the variation of Mg II line strength with impact parameter as a function of the color and mass of the foreground galaxy. The foreground sample is divided into four sub-samples in terms of their stellar masses and rest-frame colors: a high-mass blue sample is defined with $\log_{10}(M_{\text{stellar}}) > 9.88 M_\odot$, a low-mass blue sample comprises galaxies with $\log_{10}(M_{\text{stellar}}) < 9.88 M_\odot$. A high-mass red sample is defined with $\log_{10}(M_{\text{stellar}}) > 10.68 M_\odot$ and a low-mass red sample comprises galaxies with $\log_{10}(M_{\text{stellar}}) < 10.68 M_\odot$, as shown in Figure 2. These mass ranges are chosen so that each sample has approximately the same number of background spectra and so that the high-mass blue sample and the low-mass red sample have more or less the same mean stellar mass, thereby enabling a direct comparison between blue and red galaxies at the same mass. For each sample, a co-added spectrum is produced for impact parameters $b < 50$ kpc, $50$ kpc $< b < 65$ kpc, and $65$ kpc $< b < 80$ kpc.

The absorption EWs measured in the 12 co-added spectra are plotted in Figure 4 and displayed in Table 1. The error bars on mass give the $1\sigma$ spread of the foreground galaxy mass in that bin. It can clearly be seen that blue foreground galaxies are associated with stronger Mg II absorption relative to red galaxies, especially at small impact parameters. To study this in more detail, we also create a single relatively narrow intermediate bin centered in mass on the overlap region (i.e., $10.1 M_\odot > \log_{10}(M_{\text{stellar}}) > 10.6 M_\odot$, shaded region in...

![Figure 3](image-url) Co-added spectra for all foreground galaxies, at three different impact parameters: $b < 40$ kpc, $40$ kpc $\leq b < 60$ kpc, and $100$ kpc $\leq b < 200$ kpc. The composite spectra are here smoothed with a box of 5 Å and are displaced vertically for presentation.

![Figure 4](image-url) Comparison of Mg II EW for different samples of the foreground galaxies selected by their mass and rest-frame color. The blue and red points are for blue and red galaxies, respectively. Three impact parameter bins in the range of 0–50 kpc (filled circles), 50–65 kpc (diamonds), and 65–80 kpc (squares) are used. Evidently there is a strong dependence on color for low impact parameter bins. This effect is not visible as we probe further away from the galaxy. There is also a mass dependence on the blue sample which is not evident for the red sample. In the left panel, we apply a small offset of 0.1 in mass for clarity. The shaded region in both panels indicate the common mass bin chosen to compare the color dependence on Mg II absorption strength.
is associated with material entrained in outflows driven by star
region of mass, the blue galaxies have almost eight times larger
Figure 4). This is shown in the panel to the right. In this overlap
dependence within the blue sample, again most clearly seen
at smaller impact parameters. This may be present for the red
galaxies, and all red galaxies. We also distinguish between galaxies that
those not in groups. The resulting radial profiles are shown in
Figure 5 and also displayed in Table 2.

The non-group galaxies, in the left-hand panel, show a profile
that drops steeply with radius beyond about 70 kpc. To ease
comparison with previous work, we parameterize the radial
profiles in terms of a singular isothermal sphere (SIS) as in
Tinker & Chen (2008). We show later in Section 3.3 that the
Mg ii absorption is certainly not spherically distributed, and thus
our choice has no physical meaning per se, and is used simply as a
convenient parameterization. In a future paper we will explore
more physically motivated models for the radial profile.

In the isothermal sphere distribution, the distribution of Mg ii
gas is characterized as

\[
EW_{\text{r}} = \begin{cases} 
\frac{EW_0}{\sqrt{(b^2/a_b^2 + 1)}} \arctan \sqrt{\frac{R_{\text{gas}}^2 - b^2}{b^2 + a_b^2}}, & \text{if } b \leq R_{\text{gas}} \\
0, & \text{otherwise} 
\end{cases}
\]  

(2)

Here the core radius \(a_b\) is defined to be \(a_b = 0.2 R_{\text{gas}}\) and does not affect \(EW_{\text{r}}\) at large \(b\). We fit this model to the EW profiles of the non-group galaxies. We find \(R_{\text{gas}} = 115.2 \pm 2.3\) kpc (physical) and \(EW_0 = 0.7 \pm 0.12\) for "all" galaxies, \(R_{\text{gas}} = 107.6 \pm 1.3\) kpc (physical) and \(EW_0 = 1.1 \pm 0.14\) for blue galaxies, and \(R_{\text{gas}} = 118.1 \pm 5.0\) kpc (physical) and \(EW_0 = 0.3 \pm 0.11\) for red galaxies. For completeness, the results for all galaxies (group plus non-group) is very similar \(R_{\text{gas}} = 114.7 \pm 1.5\) kpc (physical) and \(EW_0 = 0.71 \pm 0.12\). It is noticeable that the \(R_{\text{gas}}\) for the blue and the red galaxies are very similar despite large changes in \(EW_0\).

Although the methodologies are different, in the sense that we measure a spatially averaged absorption, our own analysis clearly confirms the trend between impact parameter \(b\) and EW for galaxies, for example as observed in Chen et al. (2010a) with quasar absorption lines in \(z < 0.5\) SDSS galaxies. However, Chen et al. (2010a) did not find any color dependence for non-group \(z < 0.5\) galaxies. As shown in Figure 4, our co-added spectra (at \(< z > \sim 0.7\)) show a strong dependence on the color of the host galaxy for \(b \leq 70\) kpc. The reason for this

| Sample | \(b^a\) | (mass)\(^b\) | Number\(^c\) | EW\(^d\) |
|--------|-------|--------------|------------|--------|
| Blue galaxies (\(M\_\text{stellar} < 9.88 M_\odot\)) | \(b < 50\) | 9.57 ± 0.20 | 116 | 0.50 ± 0.08 |
| | 50 < \(b < 65\) | 9.57 ± 0.17 | 90 | 0.26 ± 0.06 |
| | 65 < \(b < 80\) | 9.6 ± 0.20 | 126 | 0.12 ± 0.08 |
| Blue galaxies (\(M\_\text{stellar} \geq 9.88 M_\odot\)) | \(b < 50\) | 10.25 ± 0.28 | 141 | 0.88 ± 0.07 |
| | 50 < \(b < 65\) | 10.25 ± 0.27 | 122 | 0.32 ± 0.05 |
| | 65 < \(b < 80\) | 10.28 ± 0.32 | 154 | 0.11 ± 0.04 |
| Red galaxies (\(M\_\text{stellar} < 10.68 M_\odot\)) | \(b < 50\) | 10.31 ± 0.30 | 104 | 0.11 ± 0.05 |
| | 50 < \(b < 65\) | 10.35 ± 0.25 | 75 | 0.08 ± 0.03 |
| | 65 < \(b < 80\) | 10.35 ± 0.27 | 102 | 0.09 ± 0.04 |
| Red galaxies (\(M\_\text{stellar} \geq 10.68 M_\odot\)) | \(b < 50\) | 10.94 ± 0.20 | 99 | 0.20 ± 0.07 |
| | 50 < \(b < 65\) | 10.93 ± 0.16 | 84 | 0.13 ± 0.03 |
| | 65 < \(b < 80\) | 10.97 ± 0.21 | 115 | 0.09 ± 0.03 |

Notes.
\(^a\) Range of impact parameter, in physical kpc.
\(^b\) Mean mass of the foreground galaxies co-added, in \(M\_\text{stellar}/M_\odot\).
\(^c\) Measured rest-frame equivalent width in Å.
\(^d\) Number of foreground–background pairs.

3.2. The Radial Dependence in Group and Non-group Environments

To study further the radial profile around galaxies, we divide
the sample more finely in \(b\) but lump together all blue galaxies
and all red galaxies. We also distinguish between galaxies that
are identified as lying in groups (as described in Section 2) and

Table 1

Measured Equivalent Width of the Complete Sample as a Function of Stellar Mass and Rest-frame Color

| Sample | \(b^a\) | (mass)\(^b\) | Number\(^c\) | EW\(^d\) |
|--------|-------|--------------|------------|--------|
| Blue galaxies (\(M\_\text{stellar} < 9.88 M_\odot\)) | \(b < 50\) | 9.57 ± 0.20 | 116 | 0.50 ± 0.08 |
| | 50 < \(b < 65\) | 9.57 ± 0.17 | 90 | 0.26 ± 0.06 |
| | 65 < \(b < 80\) | 9.6 ± 0.20 | 126 | 0.12 ± 0.08 |
| Blue galaxies (\(M\_\text{stellar} \geq 9.88 M_\odot\)) | \(b < 50\) | 10.25 ± 0.28 | 141 | 0.88 ± 0.07 |
| | 50 < \(b < 65\) | 10.25 ± 0.27 | 122 | 0.32 ± 0.05 |
| | 65 < \(b < 80\) | 10.28 ± 0.32 | 154 | 0.11 ± 0.04 |
| Red galaxies (\(M\_\text{stellar} < 10.68 M_\odot\)) | \(b < 50\) | 10.31 ± 0.30 | 104 | 0.11 ± 0.05 |
| | 50 < \(b < 65\) | 10.35 ± 0.25 | 75 | 0.08 ± 0.03 |
| | 65 < \(b < 80\) | 10.35 ± 0.27 | 102 | 0.09 ± 0.04 |
| Red galaxies (\(M\_\text{stellar} \geq 10.68 M_\odot\)) | \(b < 50\) | 10.94 ± 0.20 | 99 | 0.20 ± 0.07 |
| | 50 < \(b < 65\) | 10.93 ± 0.16 | 84 | 0.13 ± 0.03 |
| | 65 < \(b < 80\) | 10.97 ± 0.21 | 115 | 0.09 ± 0.03 |
The discrepancy is unclear to us. It might reflect the higher redshifts, and thus stronger star formation activity, of blue galaxies in our own sample, but the redshift difference is small enough to make this interpretation unlikely in our view. The right-hand panel in Figure 5 shows the radial profile around galaxies, which are selected to lie in groups. This is somewhat artificial, essentially because the absorption from each line of sight will contribute to several different $b$ bins, corresponding to the different members of the group. This will automatically produce a flatter distribution of Mg II with impact parameter. We include it in this paper only for comparison with previous work (e.g., Chen et al. 2010a). With this caveat in mind, we see a much flatter absorption profile that extends beyond 140 kpc, substantially larger than seen for isolated, non-group galaxies. It should be noted that the dependence on color that we have observed for non-group galaxies is also evident in group galaxies at small radii $b < 70$ kpc but is much weaker at larger radii.

A more physically motivated approach would be to determine the Mg II absorption profile around a particular location in the group, e.g., either the geometric center defined in C. Knobel et al. (2011, in preparation) as the geometric mean of the positions in the sky of all the group members, or the location of the most massive member of the group. The results are shown in Figure 6. The left panel is for group centers and the right panel is for the most massive galaxies in the groups.

In order to try to understand the profiles in the right-hand panel of Figure 5 and in Figure 6, we construct a simple model in which the absorption profile around each group galaxy is given by that which we have measured around non-group galaxies—i.e., the model assumes that the absorption of the group can be represented by a simple superposition of the absorption of individual members. Knowing the locations of each individual members within each group, we can predict, for the ensemble of group members, the average profile that would be expected if all group members exhibited the same
The Astrophysical Journal, 743:10 (11pp), 2011 December 10

Bordoloi et al.

Table 2

| Environmenta | Colorb | b | Numberd | EWc |
|--------------|--------|---|---------|-----|
| 0 < b < 40   | All galaxies | 198 | 0.48 ± 0.05 |
| 40 < b < 60  | 60 < b < 80 | 300 | 0.38 ± 0.05 |
| 80 < b < 120 | 120 < b < 160 | 475 | 0.18 ± 0.05 |
| 120 < b < 160 | 60 < b < 120 | 1273 | 0.04 ± 0.02 |
| 120 < b < 160 | 80 < b < 120 | 1785 | <0.012 |
| 120 < b < 160 | Non-group galaxies | 0 < b < 40 | 116 | 0.77 ± 0.08 |
| 40 < b < 60  | 60 < b < 80 | 159 | 0.46 ± 0.05 |
| 80 < b < 120 | 120 < b < 160 | 278 | 0.13 ± 0.06 |
| 120 < b < 160 | 60 < b < 120 | 742 | 0.02 ± 0.01 |
| 120 < b < 160 | Red galaxies | 0 < b < 40 | 82 | 0.18 ± 0.05 |
| 40 < b < 60  | 60 < b < 80 | 141 | 0.13 ± 0.02 |
| 80 < b < 120 | 120 < b < 160 | 197 | 0.16 ± 0.02 |
| 120 < b < 160 | 60 < b < 120 | 531 | 0.04 ± 0.01 |
| 120 < b < 160 | Group galaxies | 0 < b < 40 | 36 | 0.61 ± 0.13 |
| 40 < b < 60  | 60 < b < 80 | 116 | 0.46 ± 0.07 |
| 80 < b < 120 | 120 < b < 160 | 196 | 0.39 ± 0.05 |
| 120 < b < 160 | 60 < b < 120 | 524 | 0.16 ± 0.04 |
| 120 < b < 160 | 80 < b < 120 | 669 | 0.15 ± 0.03 |
| 120 < b < 160 | Red galaxies | 0 < b < 40 | 36 | 0.61 ± 0.03 |
| 40 < b < 60  | 60 < b < 80 | 92 | 0.35 ± 0.09 |
| 80 < b < 120 | 120 < b < 160 | 248 | 0.14 ± 0.04 |
| 120 < b < 160 | 60 < b < 120 | 320 | 0.11 ± 0.05 |
| 120 < b < 160 | 80 < b < 120 | 467 | <0.015 |
| 120 < b < 160 | 60 < b < 120 | 39 | 0.15 ± 0.04 |
| 40 < b < 60  | 60 < b < 80 | 54 | 0.13 ± 0.04 |
| 80 < b < 120 | 120 < b < 160 | 104 | 0.20 ± 0.02 |
| 120 < b < 160 | 60 < b < 120 | 276 | 0.09 ± 0.05 |
| 120 < b < 160 | 80 < b < 120 | 349 | 0.11 ± 0.02 |
| 120 < b < 160 | Red galaxies | 0 < b < 40 | 39 | 0.15 ± 0.04 |
| 40 < b < 60  | 60 < b < 80 | 54 | 0.13 ± 0.04 |
| 80 < b < 120 | 120 < b < 160 | 104 | 0.20 ± 0.02 |
| 120 < b < 160 | 60 < b < 120 | 276 | 0.09 ± 0.05 |
| 120 < b < 160 | 80 < b < 120 | 349 | 0.11 ± 0.02 |
| 120 < b < 160 | 60 < b < 120 | 477 | <0.074 |

Notes.

a Dividing the galaxies in terms of environment.
b Galaxies are divided as blue and red, in terms of their rest-frame colors.
c Range of impact parameter, in physical kpc.
d Number of foreground–background pairs.
e Measured rest-frame equivalent width in Å.

average profile as non-group members, as described by the SIS profiles in the left-hand panel of Figure 5. On average, only 2/3 of the $I_{AB} \leq 22.5$ galaxies were observed spectroscopically in zCOSMOS-bright. C. Knobel et al. (2011, in preparation) have developed a scheme to include in the groups, the missing members for which only a photo-$z$ is available. In the model we include along with the spectroscopic members, the photometric members with a probability of membership $p \geq 0.7$. We first ignore the color of the other group members, but then incorporate this information also into the model.

If, within some bin in impact parameter, we have a sample of $n$ group members, which are spread among some number of groups, and if the $n$th galaxy in this sample has $m$ fellow members of its particular group, then the final EW that is expected for this impact parameter is given by

$$\text{EW}_{\text{total}} = \frac{1}{n} \sum_{n} \left( \sum_{m} \text{EW}_m(b_m) \right).$$

Figure 7. Black points with filled circles are EW measurements for all group galaxies; the dotted lines are the input SIS profiles for all galaxies (black dotted line), blue galaxies (blue dotted line), and red galaxies (red dotted line), respectively. The dark shaded region gives the EW profile expected from the superposition model with a radial SIS profile for each halo as given by the black dotted line. The light shaded region is the expectation of the same model if we include the color information for individual group members and assign different EW profiles for red and blue galaxies as given by the dotted red and blue lines.

(A color version of this figure is available in the online journal.)

where $\text{EW}_m(b_m)$ is the EW appropriate for the SIS sphere of the galaxy type for galaxy $m$, i.e., blue or red, at the impact parameter $b_m$ that the background galaxy has to this $n$th member galaxy. In summing the EWs, we assume that the velocities associated with each galaxy are large enough that there are no saturation effects. Typical velocity dispersions within the groups in question are typically $200 \text{ km s}^{-1} < \sigma < 600 \text{ km s}^{-1}$ (Knobel et al. 2009).

The results of this exercise are shown in Figure 7 and in Figure 6. Figure 7 re-plots the average group member profile of Figure 5 but no longer differentiating in color of the foreground galaxies. Both in Figure 7 and in Figure 6, the dotted lines show what would be expected from the SIS model around a single galaxy, i.e., ignoring other group members entirely, the dark shaded area includes the effects of other group members using the superposition approach, while the light shaded area also incorporates the information on the colors of those other members. Both of the latter shaded areas provide a better representation of the extended profile observed in group galaxies. This simple superposition model provides a reasonable representation of the absorption profiles around group galaxies. Interestingly, such a model would predict that the absorption profile at $b > 100$ kpc should be independent of the color of the member galaxy, since at these radii it would be dominated by other group members. This may indeed be seen in the right-hand panel of Figure 5.

This superposition model reproduces both the extended distribution and the magnitude of the Mg II absorption profile very well using the geometrical group centers. However, while using the most massive galaxies, the result is less satisfactory. The differences between the two panels of Figure 6 arise from differences in both the models and in the observational data points, especially at large impact parameters. The difference in the models arises because the group centers are, by construction, more or less equidistant from the group members, maximizing the number which are close enough to contribute to the expected EW. Within the data, there are about 2.5 times more group
members within 100 kpc of the geometrical center as there are within 100 kpc of the most massive galaxies. The differences in the observational points are not of large significance compared with the statistical error bars. The offsets between the geometrical centers and the most massive galaxies in each group extend up to 140 kpc, producing substantial scrambling of the data in impact parameter between the two approaches. Therefore, the error bars between the plots should be largely independent.

The success of this superposition model, especially for geometrical group centers is quite surprising as it suggests that the existence of Mg ii absorption halos around galaxies may not be significantly affected by the group environment. The fact that the specific star formation rates of star-forming galaxies do not appear to depend on environment, even though the fraction of galaxies that are star-forming does (Peng et al. 2010), suggests that the source of star-formation-driven winds is likely the same in the two environments. However, tidal effects have been invoked to enhance absorption for low-redshift groups (Kacprzak et al. 2010b). One can also imagine scenarios whereby different mechanisms may counteract each other, and the accuracy of the agreement anyway allows for some non-negligible differences.

3.3. The Azimuthal Absorption Profile Around Disk Galaxies

With the current data set, we can probe the strength of Mg ii absorption lines around disk galaxies as a function of the azimuthal angle relative to the projected disk axis. As described in Section 2, we select a set of disk-dominated galaxies that lie within 40° of being edge-on. We then compute the azimuthal angle φ between the projected semi-minor axis of the disk, and the projected vector from the center of the foreground galaxy to the background galaxy. In other words, values of |φ| > 45° represent lines of sight that pass near to the plane of the disk, while values of |φ| < 45° are associated with lines of sight passing close to the symmetry (rotation) axis of the disk. If the Mg ii absorption around disk galaxies is associated with a bi-polar outflow along the disk axis, then the latter might be expected to show stronger absorption. Conversely, if the absorption was due to an extension of the disk itself, then the former would be expected to be stronger.

Figure 8 shows the radial Mg ii absorption profiles for these disk galaxies for two azimuthal bins split at |φ| = 45°. At small impact parameters, i.e., b < 50 kpc, the absorption along the disk axis is significantly stronger than in the plane of the disk. The difference disappears at larger radii because the polar quadrants have a much steeper radial decline. It is quite noticeable how flat the radial profile in the plane of the disk is. For low impact parameters, stronger Mg ii systems are observed close to the axis of the disk as compared to the plane of the disk. This effect diminishes as higher impact parameters are probed and is no longer distinguishable for very high impact parameters. The error bars in impact parameter give the standard deviation of the distribution of impact parameters within that bin.

(A color version of this figure is available in the online journal.)

Table 3

Radial Profile of Mg ii Absorption Around Disk Galaxies at Different Azimuthal Angles

| φ°   | b kpc | Number | EW (Å) |
|------|-------|--------|--------|
| 0° < | 0 < b | 35     | 23     | 1.17 ± 0.25 |
| 35 < | b < 50| 31     | 0.92 ± 0.25 |
| 50 < | b < 65| 42     | 0.47 ± 0.12 |
| 65 < | b < 80| 91     | 0.19 ± 0.10 |

| 45° < | | 60 < b | | 90° | |
|-------| | 0 < b | 35 | 25 | 0.41 ± 0.20 |
| | 35 < b | 50 | 28 | 0.47 ± 0.14 |
| | 50 < b | 65 | 49 | 0.43 ± 0.10 |
| | 65 < b | 80 | 93 | 0.24 ± 0.10 |

Notes.

a Range of impact parameter, in physical kpc.

b Range of azimuthal angles considered, in degrees

c Number of foreground–background pairs.

d Measured rest-frame equivalent width in Å.
above offer a much stronger and more robust diagnostic of this, which may account for some of the apparent differences. However, Kacprzak et al.’s (2011) result primarily refers to weaker Mg\textsc{ii} systems at relatively large impact parameters and this may account for some of the apparent differences.

Although we believe that the azimuthal effects explored above offer a much stronger and more robust diagnostic of the geometry of the Mg\textsc{ii} distribution than inclination effects, we present in this section, for completeness, the variation of Mg\textsc{ii} absorption with the inclination of the associated disk galaxies. As in the previous section, we select the foreground galaxies to be disk dominated, as defined in Section 2. We divide the disk galaxies into three bins in inclination angles, within $0^\circ < \phi < 50^\circ$, $50^\circ < \phi < 65^\circ$, and $65^\circ < \phi < 90^\circ$, for each of three bins in impact parameter, $b < 40$ kpc (black circle), $40$ kpc $< b < 60$ kpc (blue diamonds), and $60$ kpc $< b < 80$ kpc (red squares). The error bars in angle are the errors on the mean of the inclination angles within that bin.

(A color version of this figure is available in the online journal.)

Figure 10. Variation of Mg\textsc{ii} EW with inclination of disk-dominated galaxies. The three inclination angle bins with $0^\circ < \phi < 50^\circ$, $50^\circ < \phi < 65^\circ$, and $65^\circ < \phi < 90^\circ$ are used for each of three bins in impact parameter, $b < 40$ kpc (black circle), $40$ kpc $< b < 60$ kpc (blue diamonds), and $60$ kpc $< b < 80$ kpc (red squares). The error bars in angle are the errors on the mean of the inclination angles within that bin.

The main results of this study are as follows.

3.4. The Dependence of Mg\textsc{ii} Absorption on the Apparent Inclination Angle of Disk Galaxies

After the submission of our original manuscript, Kacprzak et al. (2011) have presented evidence from analysis of the inclinations of Mg\textsc{ii} selected galaxies that the distribution of Mg\textsc{ii} is co-planar with galaxy disks, especially for weak absorption systems. At first sight, this is in direct contrast to the evidence in the previous section that Mg\textsc{ii} absorption is strongest in bi-polar regions aligned with the poles of the disks. However, Kacprzak et al.’s (2011) result primarily refers to weaker Mg\textsc{ii} systems at relatively large impact parameters and this may account for some of the apparent differences.

Although we believe that the azimuthal effects explored above offer a much stronger and more robust diagnostic of the geometry of the Mg\textsc{ii} distribution than inclination effects, we present in this section, for completeness, the variation of Mg\textsc{ii} absorption with the inclination of the associated disk galaxies. As in the previous section, we select the foreground galaxies to be disk dominated, as defined in Section 2. We divide the disk galaxies into three bins in inclination angles, within $0^\circ < \phi < 50^\circ$, $50^\circ < \phi < 65^\circ$, and $65^\circ < \phi < 90^\circ$.

The sense of the inclination angles is that $i = 90^\circ$ represents an edge-on system and $i = 0^\circ$ represents a face-on system. The sample is not sub-divided in terms of azimuthal angles, i.e., for any range of inclination the derived EW are averaged over all azimuthal angles. Figure 10 (Table 5) shows the variation of Mg\textsc{ii} absorption around disk galaxies as a function of inclination for three impact parameter bins. The error bars in inclination angles are the standard errors on the mean of the inclination angles within that bin. We find no significant trends of EW with the apparent inclination of disk-dominated galaxies.

We plan to present a more comprehensive exploration of the dependence of EW on both inclination and azimuthal angles for a range of spatial distributions in a future paper. In this we will explore any differences between the two approaches described by Kacprzak et al. (2011) and here.

4. CONCLUSIONS

In this work, we have mapped the average spatial distributions of Mg\textsc{ii} gas around a set of zCOSMOS-bright galaxies at $0.5 < z < 0.9$, using co-added spectra from a large sample of background zCOSMOS-deep objects. We divide the foreground galaxies in terms of their rest-frame color and mass, and also by their location in or out of groups. We also constructed average profiles for the groups around both their geometric centers and around the most massive galaxies in each group. Finally, we examined the azimuthal dependence of Mg\textsc{ii} absorption around inclined disk galaxies and investigated the dependence of Mg\textsc{ii} absorption on apparent inclination of disk-dominated galaxies. The main results of this study are as follows.

1. We find that blue foreground galaxies are associated with much stronger Mg\textsc{ii} absorption compared to red galaxies,
particularly at small impact parameters. Within the overlapping mass range (and for $b < 50$ kpc), blue galaxies are associated with absorption systems which are about eight times stronger in EW than red galaxies. This is consistent with the idea that entrained material in outflows driven by star formation are responsible for the Mg II absorption, but could also indicate inflowing material feeding star formation.

2. There is also a clear correlation between Mg II absorption line strength and the host galaxy stellar mass for the blue galaxies, especially at low impact parameters ($b < 50$ kpc). This may indicate a correlation with star formation rate, since the specific star formation rate in star-forming galaxies is known to be roughly constant.

3. For isolated (non-group) galaxies, the Mg II radial profile can be well represented by an SIS model similar to Chen et al. (2010a). We estimate the effective gas radius is $R_{\text{gas}} = 115.2 \pm 2.3$ kpc (physical) and $EW_0 = 0.7 \pm 0.12$. For blue galaxies it is $R_{\text{gas}} = 107.6 \pm 1.3$ kpc (physical) and $EW_0 = 1.1 \pm 0.14$, while for red galaxies it has a very similar size $R_{\text{gas}} = 118.1 \pm 5.$ kpc (physical) but much lower $EW_0 = 0.3 \pm 0.11$.

4. Galaxies in groups have a much flatter absorption profile that extends beyond 140 kpc. The color dependence on Mg II line strength is maintained for the group galaxies, especially at small $b$. Likewise, the average Mg II absorption profiles constructed around a specific point in each group, either a geometric center or the location of the most massive member of the group, are also substantially more extended than the profiles around isolated galaxies.

5. The extended distributions of Mg II absorption in groups can however be well reproduced with a simple model that superposes the absorption profiles of individual group members, assuming each of these is represented by an absorption profile that is the same as for isolated galaxies. This might indicate that the transport of Mg II out of galaxies may not be strongly affected (within our uncertainties) by the intra-group medium, or by other processes that may occur in groups, such as tidal interactions.

6. For disks that are seen close to edge-on (inclinations greater than 50° in our convention), the Mg II absorption strength is about three times stronger in regions above the projected rotation axis of the disks (i.e., along the apparent minor axis) than in the plane of the disks (i.e., along the apparent major axis) for small impact parameters ($b < 40$ kpc). We suggest that this is a strong evidence that Mg II absorption is associated with bipolar regions aligned with the disk axis, presumably indicating the presence of bipolar Galactic winds. The azimuthal asymmetry is not seen at larger distances from the galaxies, indicating that the radial fall-off in absorption is faster along the disk axis than in the plane of the disk.

7. We find no statistically significant correlation between inclinations of disk galaxies and the (azimuthally averaged) Mg II absorption strength.

We thank Nicolas Bouché for helpful discussions on this work. We also thank the anonymous referee for providing constructive and insightful comments that helped in improving the paper. This work has been supported by the Swiss National Science Foundation and is based on observations undertaken at the European Southern Observatory (ESO) Very Large Telescope (VLT) under Large Program 175.A-0839.

REFERENCES

Bergeron, J., & Stasińska, G. 1986, A&A, 169, 1
Bernet, M. L., Miniati, F., Lilly, S. J., Kronberg, P. P., & Dessauges-Zavadsky, M. 2008, Nature, 454, 302
Bolzonella, M., Kovač, K., Pozzetti, L., et al. 2010, A&A, 524, A76
Bolzonella, M., Miralles, J.-M., & Pelló, R. 2000, A&A, 363, 476
Bouché, N. 2008, MNRAS, 389, L18
Bouché, N., Murphy, M. T., Péroux, C., Csabai, I., & Wild, V. 2006, MNRAS, 371, 495
Bowen, D. V., Blades, J. C., & Pettini, M. 1995, ApJ, 448, 662
Charlton, J. C., & Churchill, C. W. 1998, ApJ, 499, 181
Charlton, J. C., Ding, J., Zonak, S. G., et al. 2003, ApJ, 589, 111
Chelouche, D., & Bowen, D. V. 2010, ApJ, 722, 1821
Chen, H., Helsby, J. E., Gauthier, J., et al. 2010a, ApJ, 714, 1521
Chen, H., & Tinker, J. L. 2008, ApJ, 687, 745
Chen, H., Wild, V., Tinker, J. L., et al. 2010b, ApJ, 724, L176
Churchill, C., & Charlton, J. 1999, BAAS, 31, 1451
Churchill, C., Steidel, C., & Kacprzak, G. 2005a, in ASP Conf. Ser. 311, Extra-Planar Gas, ed. R. Braun (San Francisco, CA: ASP), 387
Churchill, C. W., Kacprzak, G. G., & Steidel, C. C. 2005b, in Proc. IAU Colloq. 199, Probing Galaxies through QSO Absorption Lines, ed. P. R. Williams, C.-G. Shu, & B. Menard (Cambridge: Cambridge Univ. Press), 24
Churchill, C. W., Mellon, R. R., Charlton, J. C., et al. 2000, ApJS, 130, 91
Cibinel, O., Capak, P., Salvato, M., et al. 2009, ApJ, 699, 1236
Kacprzak, G. G., Churchill, C. W., Ceverino, D., et al. 2010a, ApJ, 711, 533
Kacprzak, G. G., Churchill, C. W., Evans, J. L., Murphy, M. T., & Steidel, C. C. 2011, MNRAS, in press (arXiv:1106.3065)
Kacprzak, G. G., Churchill, C. W., Steidel, C. C., & Murphy, M. T. 2008, AJ, 135, 922
Kacprzak, G. G., Murphy, M. T., & Churchill, C. W. 2010b, MNRAS, 406, 445
Knobel, C., Lilly, S. J., Iovino, A., et al. 2009, ApJ, 697, 1842
Koekemoer, A. M., Aussel, H., Calzetti, D., et al. 2007, ApJ, 172, 196
Lanzetta, K. M., Turnshek, D. A., & Wolfe, A. M. 1987, ApJ, 322, 739
Lilly, S. J., Le Brun, V., Maier, Ch., et al. 2009, ApJ, 184, 218
Lilly, S. J., Le Fèvre, O., Renzini, A., et al. 2007, ApJ, 172, 70
Ménard, B., Wild, V., Nestor, D., Quider, A., & Zibetti, S. 2009, MNRAS, submitted (arXiv:0912.3263)
Nestor, D. B., Johnson, B. D., Wild, V., et al. 2010, MNRAS, 412, 1559
Nestor, D. B., Turnshek, D. A., & Rao, S. M. 2005, ApJ, 628, 637
Nestor, D. B., Turnshek, D. A., Rao, S. M., & Quider, A. M. 2007, ApJ, 658, 185
Peng, Y., Lilly, S. J., Kovač, K., et al. 2010, ApJ, 721, 193
Petitjean, P., & Bergeron, J. 1990, A&A, 231, 309
Prochter, G. E., Prochaska, J. X., & Burles, S. 2006, ApJ, 639, 766
Rao, S. M., Turnshek, D. A., & Nestor, D. B. 2006, ApJ, 636, 610
Rigby, J. R., Charlton, J. C., & Churchill, C. W. 2005, ApJ, 656, 743
Rubin, K. H. R., Weiner, B. J., Koo, D. C., et al. 2010, ApJ, 719, 1503
Sargent, W. L. W., Steidel, C. C., & Boksenberg, A. 1988, ApJ, 334, 22
Scarlata, C., Carollo, C. M., Lilly, S., et al. 2007, ApJ, 72, 406
Scoville, N., Aussel, H., Brusa, M., et al. 2007, ApJS, 172, 1
Steidel, C. C. 1995, in Proc. ESO Workshop, QSO Absorption Lines, ed. G. Meylan (Berlin: Springer), 139
Steidel, C. C., Erb, D. K., Shapley, A. E., et al. 2010, ApJ, 717, 289
Steidel, C. C., & Sargent, W. L. W. 1992, ApJS, 80, 1
Tinker, J. L., & Chen, H. 2008, ApJ, 679, 1218
Weiner, B. J., Coil, A. L., Prochaska, J. X., et al. 2009, ApJ, 692, 187
Zibetti, S., Ménard, B., Nestor, D. B., et al. 2007, ApJ, 658, 161
Zucca, E., Bardelli, S., Bolzonella, M., et al. 2009, A&A, 508, 1217