A New Compact Alternating Direction Implicit Method for Solving Two Dimensional Time Fractional Diffusion Equation With Caputo-Fabrizio Derivative

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Abstract. In this paper, a new compact alternating direction implicit (ADI) difference scheme is proposed for the solution of two dimensional time fractional diffusion equation. Theoretical considerations are discussed. We show that the proposed method is fourth order accurate in space and two order accurate in time. The stability and convergence of the compact ADI method are presented by the Fourier analysis method. Numerical examples confirm the theoretical results and high accuracy of the proposed scheme.

1. Introduction

Fractional differential equations (FDEs), which deal with derivatives and integrals of any arbitrary real or complex order, have been highly regarded by researchers. Fractional calculus has a history of more than 300 years, yet its applicability in different domains has been realized recently. In the last three decades, the subject witnessed a significant growth of research [1, 2]. Recently, fractional differential equations have been widely used in various fields such as muscular blood vessel modeling [3], non-linear oscillation of earthquakes [4], control theory [5], financial economics [6], biotechnology [7] and etc [8–11]. Fractional derivatives model the various dynamical processes and they carry information regarding their present as well as past states. In order to characterize memory property of complex systems, one needs to employ the non-integer order derivatives because these operators give a complete description of different physical processes with dissipation and long-range interaction [12]. There are many numerical methods proposed for solving the FDEs up to now, e.g., finite difference method [13], spectral methods [14–16], finite element method [17, 18], RBF [19] and etc. Low-order finite difference schemes are not accurate enough for solving many problems in science. Recently the focus has shifted to high order compact finite difference methods [20–23]. The advantage of the high order compact finite difference is that they give high accuracy on small stencils with greater computational efficiency [24]. The original split type method was introduced by Peaceman and Rachford in 1955. Their method is called the alternating direction implicit method. Such schemes split multi-dimensional problems to a series of one-dimensional ones which are much easier to solve [25].
Recently in [28], and a compact ADI method for solving two-dimensional Riesz space fractional diffusion was given recently in [28], and a compact ADI method for solving two-dimensional Riesz space fractional diffusion equation proposed in [29].

Caputo and Fabrizio have suggested a new definition of a fractional derivative which we will use in this paper [30]. In this paper, we consider the following two-dimensional time fractional diffusion equation describing sub-diffusive phenomena with a non-homogeneous term

$$ \frac{\mathcal{C}D^\gamma_t u(x, y, t)}{0} = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(x, y, t), \quad (x, y) \in \Omega, \quad 0 < t \leq T, $$

with the initial condition

$$ u(x, y, 0) = w(x, y), \quad (x, y) \in \Omega, $$

and the Dirichlet boundary conditions

$$ u(L_1, y, t) = \varphi_1(y, t), u(L_2, y, t) = \varphi_2(y, t), $$

$$ u(x, L_3, t) = \psi_1(x, t), u(x, L_4, t) = \psi_2(x, t), \quad t \geq 0, $$

where $\Omega = (L_1, L_2) \times (L_3, L_4)$. Here $0 < \gamma < 1$ and $\mathcal{C}D^\gamma_t u$ denotes the Caputo-Fabrizio fractional derivative of the function $u(x, y, t)$ defined as:

$$ \mathcal{C}D^\gamma_t u(x, y, t) = \frac{M(\gamma)}{1 - \gamma} \int_0^t \frac{\partial u(x, y, s)}{\partial s} e^{\sigma(t-s)} ds, $$

where $M(\gamma)$ is a normalization function such that $M(0) = M(1) = 1$ and $\sigma = \frac{T}{1 - \gamma}$. Our goal is to present a compact ADI scheme to solve (1)-(3) based on the new fractional derivative. The unconditionally stable result is derived.

The paper is organized as follows. In Section 2, we give a compact ADI difference scheme for two-dimensional time fractional diffusion equation. Then in Section 3, we present the analysis of stability and convergence for the presented scheme. In Section 4, some numerical results using the fourth order compact finite difference scheme are carried out. Finally, this paper ends with conclusions in Section 5.

2. Compact ADI scheme for 2D time-fractional diffusion equation with Caputo-Fabrizio derivative

The domain $\Omega \times [0, T]$ of (1)–(3) is divided into a uniform grid of mesh points $(x_j, y_k, t_n)$ with $x_j = L_1 + jh_x$, $j = 0, 1, \ldots, N_x + 1$, $y_k = kh_y$, $k = 0, 1, \ldots, N_y + 1$, and $t_n = nt$, $n = 0, 1, \ldots, N$. Here $N_x$, $N_y$ and $N$ are positive integers, $h_x = (L_2 - L_1)/(N_x + 1)$ and $h_y = (L_4 - L_3)/(N_y + 1)$ are the mesh-widths in $x, y$, respectively, with $h_x/h_y$ bounded from below and above, and $\tau = T/N$ is the time step.

A discrete approximation to the $\mathcal{C}D^\gamma_t u(x, y, t)$ at $(x_j, y_k, t_n)$ can be obtained by the following approxima-
Consider the partial differential equation (6) gives

\[ \frac{\partial u}{\partial x}(x, y, s) = \frac{\partial u}{\partial x}(x, y, s) + \frac{\partial u}{\partial y}(x, y, s) = f(x, y, t). \]

A fourth-order compact finite difference scheme for this equation is given as

\[ \left( \frac{\delta^2_x u^n_{jk} + \delta^2_x u^n_{jk}}{h^2_x} + \frac{\delta^2_y u^n_{jk} + \delta^2_y u^n_{jk}}{h^2_y} \right) = \left( \frac{\delta^4_x f^n_{jk} + \delta^4_y f^n_{jk}}{12} \right) + O(h^4_x + h^4_y). \]

Proof. Based on the definition of operator \( \delta_y \), we have the following relation for (6) at point \( (x_j, y_k, t_n) \)

\[ \frac{\delta^2_x u^n_{jk} + \delta^2_y u^n_{jk} - \delta^n_{jk}}{h^2_x} f^n_{jk}, \]

where

\[ \delta^n_{jk} = \frac{h^2_x}{12} \left( \frac{\partial^4 u^n}{\partial x^4} \right)_{jk} + \frac{h^2_y}{12} \left( \frac{\partial^4 u^n}{\partial y^4} \right)_{jk} + O(h^4_x + h^4_y). \]

In order to obtain a fourth-order scheme, the fourth derivatives of \( u \) in Eq. (9) should be approximated. Eq. (6) gives

\[ \frac{\partial^4 u^n}{\partial x^4} |_{jk} = \left( \frac{\delta^2 f}{\delta x^2} - \frac{\delta^2 f}{\delta x^2 \delta y} \right) |_{jk} = \delta^2_{x} f^n_{jk} - \delta^2_{x} \delta^2_{y} u^n_{jk}, \]

\[ \frac{\partial^4 u^n}{\partial y^4} |_{jk} = \left( \frac{\delta^2 f}{\delta y^2} - \frac{\delta^2 f}{\delta y^2 \delta x} \right) |_{jk} = \delta^2_{y} f^n_{jk} - \delta^2_{y} \delta^2_{x} u^n_{jk}. \]
By substituting Eqs. (10) and (11) in Eq. (9) we get
\[ \tau_{jk}^n = \frac{1}{12} (\delta_{x,j}^2 f_{jk} + \delta_{y,j}^2 f_{jk}) - \frac{1}{12} \left( \frac{\partial^2}{\partial y^2} u_{jk}^n + \frac{\partial^2}{\partial x^2} u_{jk}^n \right) + O(h_x^4 + h_y^4), \] (12)
hence by substituting Eq. (12) in Eq. (8) the lemma is proved. \( \square \)

**Lemma 2.2.** Suppose \( \psi(t) \in C^2[0, t_n] \). Let
\[ A = \sum_{i=1}^{n} \int_{b_{i-1}}^{b_i} \left[ \frac{\psi'(t) - \psi(t_{i-1})}{\Delta t} \right] e^{-\sigma(t_{i-1})} dt, \]
then
\[ |A| \leq \frac{(\Delta t)^2}{\sigma} \max_{0 \leq t \leq h_n} |\psi''(t)| \left( \frac{\sigma^2}{12} - \frac{\sigma^3 \Delta t}{24} + \cdots \right) t_n. \]

**Proof.** Using the Taylor series expansion with integral remainder, we have
\[ A = \frac{1}{\Delta t} \sum_{i=1}^{n} \int_{b_{i-1}}^{b_i} \left[ \psi''(s - t_{i-1}) ds - \int_{b_{i-1}}^{b_i} \psi''(s) (t_i - s) ds \right] e^{-\sigma(t_{i-1})} dt \]
\[ = \frac{1}{\Delta t} \sum_{i=1}^{n} \int_{b_{i-1}}^{b_i} \left[ (s - t_{i-1}) \left( e^{-\sigma(t_{i-1})} - e^{-\sigma(t_i)} \right) - (t_i - s) \left( e^{-\sigma(t_i)} - e^{-\sigma(t_{i-1})} \right) \right] \psi''(s) ds. \]
Since,
\[ \int_{b_{i-1}}^{b_i} \left[ (s - t_{i-1}) \left( e^{-\sigma(t_{i-1})} - e^{-\sigma(t_i)} \right) - (t_i - s) \left( e^{-\sigma(t_i)} - e^{-\sigma(t_{i-1})} \right) \right] ds \]
\[ = (\Delta t)^4 e^{-\sigma(\Delta t(n-l))} \left( \frac{\sigma^2}{12} - \frac{\sigma^3 \Delta t}{24} + \cdots \right) \]
\[ \leq (\Delta t)^4 \left( \frac{\sigma^2}{12} - \frac{\sigma^3 \Delta t}{24} + \cdots \right), \]
and \( e^{-\sigma(\Delta t(n-l))} < 1 \) for \( l = 1, 2, \ldots, n \) and \( t_n = n \Delta t \), hence the result will be achieved. \( \square \)

Now, using the Lemma (2.2) we get
\[ M(y) \int_{b_{i-1}}^{b_i} \left( \frac{\partial u}{\partial s} \bigg|_{(x,y,s)} - \frac{u_{jk} - u_{jk-1}}{\tau} \right) e^{-\sigma(t_{i-1})} ds \]
\[ \leq \frac{M(y)}{1 - \gamma} \frac{\tau^2}{\sigma} \max_{0 \leq t \leq h_n} |\psi''(t)| \left( \frac{\sigma^2}{12} - \frac{\sigma^3 \tau}{24} + \cdots \right) \sum_{i=1}^{n} e^{-\sigma(t_{i-1})} = O(\tau^2), \]
Let
\[ \tilde{D} \psi u(x, y, t) = \frac{M(y)}{1 - \gamma} \int_{0}^{t} \frac{\partial u(x, y, s)}{\partial s} e^{-\sigma(t-s)} ds = g(x, y, t), \] (13)
so Eq. (1) can be written as follows
\[ \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = g(x, y, t) - f(x, y, t). \] (14)
Using lemma (2.1), we can write Eq. (14) at the point \((x_j, y_k, t_n)\) as follows

\[
\left( \frac{\delta_x^2}{h_x^2} + \frac{1}{12 h_x^2} \frac{\partial^2}{\partial y^2} \right) u^n_{jk} = \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) \left( f^n_{jk} - f^0_{jk} \right) + O(h^4).
\]

Therefore, from Eq. (5) and lemma (2.2) we have

\[
\left( \frac{\delta_x^2}{h_x^2} + \frac{1}{12 h_x^2} \frac{\partial^2}{\partial y^2} \right) u^n_{jk} = \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) \times \\
\left[ \frac{1}{12 t} \left( d_{n,n} u^n_{jk} - \sum_{i=1}^{n-1} (d_{i,n} - d_{n,i}) u'_{jk} + d_{1,n} u^0_{jk} \right) \right] - \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) f^n_{jk} + O(h^4 + \tau^2),
\]

where \(M(y) = 1\). Let \(\nu = \frac{d_{n,n}}{12 t}\) then

\[
\left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) u^n_{jk} = \frac{1}{12 n} \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) \times \\
\left[ \sum_{i=1}^{n-1} (d_{i,n} - d_{n,i}) u'_{jk} + d_{1,n} u^0_{jk} \right] - \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) f^n_{jk} + O(h^4 + \tau^2).
\]

Using factorization, we can rewrite Eq. (17) as

\[
\left( 1 + \frac{1}{12 t} - \frac{\delta_x^2}{h_x^2} \right) \left( 1 + \frac{1}{12 t} - \frac{\delta_y^2}{h_y^2} \right) u^n_{jk} = \frac{1}{144} + \frac{\delta_x^2}{12 t} \left[ \sum_{i=1}^{n-1} (d_{i,n} - d_{n,i}) u'_{jk} + d_{1,n} u^0_{jk} \right] \\
- \frac{1}{\nu} \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) f^n_{jk} + O(h^4 + \tau^2).
\]

To give a numerical scheme we need to approximate the term \(\left( \frac{1}{12 t} - \frac{\delta_x^2}{h_x^2} \right) u^n_{jk}\) on the right hand side Eq. (18) and omit the last term. As in paper [31] we use the following approximation

\[
u^n_{jk} = \begin{cases} 
2 u^n_{jk} - u^{n-2}_{jk}, & n > 1, \\
0, & n = 1.
\end{cases}
\]

Using the Taylor series expansion we have

\[
u(x_j, y_k, t_{n-1}) = u(x_j, y_k, t_n) + \tau u'(x_j, y_k, t_n) + \frac{\tau^2}{2} u''(x_j, y_k, t_n) + O(\tau^3),
\]

\[
u(x_j, y_k, t_{n-2}) = u(x_j, y_k, t_n) + 2 \tau u'(x_j, y_k, t_n) + 2 \tau^2 u''(x_j, y_k, t_n) + O(\tau^3).
\]

So \(u^n_{jk} = 2 u^{n-1}_{jk} - u^{n-2}_{jk} + O(\tau^2)\). Define \(\mu_x = \frac{1}{12 t}\) and \(\mu_y = \frac{1}{12 t}\) for \(n = 1\), we have

\[
\left( 1 + \frac{\delta_x^2}{12 t} - \mu_x^2 \right) \left( 1 + \frac{\delta_y^2}{12 t} - \mu_y^2 \right) u^n_{jk} = \left( \frac{1}{144} + \mu_x \mu_y \right) u^0_{jk} \\
+ \frac{\delta_x^2}{12 t} + \frac{\delta_y^2}{12 t} + 1 \right) f^n_{jk} + O(h^4 + \tau^2).
\]
and for \( n \geq 2 \), we obtain

\[
\left( 1 + \left( \frac{1}{12} - \mu_x \right) \delta_y^2 \right) \left( 1 + \left( \frac{1}{12} - \mu_y \right) \delta_y^2 \right) u_{jk}^n
\]

\[
= \left( \frac{1}{144} + \mu_x \mu_y \right) \left( u_{jk}^{n-1} - u_{jk}^{n-2} \right) + \frac{1}{d_{jk}} \left( \delta_y^2 + \delta_x^2 + 1 \right) u_{jk}^0 + \frac{1}{12} \delta_y^2 + \delta_x^2 + 1 \right) \sum_{l=1}^{n-1} (d_{l+1,n} - d_{l,n}) u_{lk}^l + d_{1,n} u_{jk}^0 \]

\[
+ \frac{1}{6} \left( \delta_y^2 + \delta_x^2 + 1 \right) f_{jk}^n.
\]

(20)

Therefore, for \( n = 1 \) and \( n \geq 2 \) we obtain the following compact ADI scheme that is implemented in two steps

\[
\begin{align*}
\left( 1 + \left( \frac{1}{12} - \mu_x \right) \delta_y^2 \right) U_{jk}^n
\left( 1 + \left( \frac{1}{12} - \mu_y \right) \delta_y^2 \right) U_{jk}^n
\end{align*}
\]

\[
= \left( \frac{1}{144} + \mu_x \mu_y \right) \left( u_{jk}^{n-1} - u_{jk}^{n-2} \right) + \frac{1}{d_{jk}} \left( \delta_y^2 + \delta_x^2 + 1 \right) u_{jk}^0 + \frac{1}{12} \delta_y^2 + \delta_x^2 + 1 \right) \sum_{l=1}^{n-1} (d_{l+1,n} - d_{l,n}) u_{lk}^l + d_{1,n} u_{jk}^0 \]

\[
+ \frac{1}{6} \left( \delta_y^2 + \delta_x^2 + 1 \right) f_{jk}^n.
\]

(21)

Let the (known) right-hand side of Eq. (21) be denoted by \( f \), then

\[
U_{jk}^n + \left( \frac{1}{12} - \mu_x \right) \left( U_{jk+1}^n - 2U_{jk}^n + U_{jk-1}^n \right) = f_{jk}^n
\]

(22)

Thus, for each \( 1 \leq k \leq N_x, n = 1, 2, \ldots, N_y \), a simple tridiagonal system is solved for \( U_{jk}^n \).

After calculating \( U_{jk}^n \), we obtain \( U_{jk}^n \) from

\[
U_{jk}^n + \left( \frac{1}{12} - \mu_y \right) \left( U_{jk+1}^n - 2U_{jk}^n + U_{jk-1}^n \right) = U_{jk}^n
\]

For each \( 1 \leq j \leq N_y, n = 1, 2, \ldots, N_x \), a tridiagonal system of equations is solved for \( U_{jk}^n \).

There is an important point that needs to be addressed with regard to the solution of the system (22). When \( j = 1 \) or \( j = N_y \), boundary values for \( U_{jk}^n \) are required in the form of \( U_{0,k}^n \) or \( U_{N_x+1,k}^n \). We can obtain the required boundary conditions for \( U_{jk}^n \) from the second and fourth equations in (21). For example, at the \( j = 1 \) we get

\[
U_{0,k}^n = \left( 1 + \left( \frac{1}{12} - \mu_y \right) \delta_y^2 \right) U_{1,k}^n = \left( 1 + \left( \frac{1}{12} - \mu_y \right) \delta_y^2 \right) \phi_1(y_k, t_n).
\]

Theorem 2.3. The coefficient matrices in the scheme (21) is invertible.

Proof. It is clear that the coefficient matrices for (21) are strictly diagonally dominant. Thus these matrices are invertible. □
3. Error analysis

3.1. Stability analysis

To study the stability analysis of the proposed scheme, we use the Fourier method. This method is a very useful technique for analyzing the stability of a finite difference method. It uses mutually orthogonal vectors that form a basis for \( n \)-dimensional space. Thus the error term can be expanded as a linear combination of these basis vectors. Let \( \bar{U}_{jk}^n \) be the approximate solution of the scheme, and define

\[
\zeta^n = (\zeta_{1,1}^n, \zeta_{1,2}^n, \ldots, \zeta_{N_x,N_y}^n)^T.
\]

Thus we have

\[
(1 + \left( \frac{1}{12} - \mu_x \right) \delta_x^2 + \frac{1}{12} - \mu_y) \zeta_{jk}^n
= \left( \frac{1}{144} + \mu_x \mu_y \right) \delta_x^2 \delta_y^2 \zeta_{jk}^n + \left( \frac{\delta_x^2}{12} + \frac{\delta_y^2}{12} + 1 \right) \zeta_{jk}^0,
\]

and for \( n \geq 2, \)

\[
\left( 1 + \left( \frac{1}{12} - \mu_x \right) \delta_x^2 \right) \left( 1 + \left( \frac{1}{12} - \mu_y \right) \delta_y^2 \right)^n \zeta_{jk}^n
= \left( \frac{1}{144} + \mu_x \mu_y \right) \delta_x^2 \delta_y^2 (2 \zeta_{jk}^{n-1} - \zeta_{jk}^{n-2}) + \frac{1}{d_{n+1}} \left( \delta_x^2 + \delta_y^2 + 1 \right) \sum_{j=1}^{N_x} (\delta_{i+1} - 2\delta_i + \delta_{i-1}) \zeta_{ij}^n + d_{n+1} \zeta_{ij}^0.
\]

The Fourier series for \( \zeta^n(x,y) \) is

\[
\zeta^n(x,y) = \sum_{m_1=-\infty}^{\infty} \sum_{m_2=-\infty}^{\infty} \rho_n(m_1, m_2) e^{2\pi i \left( \frac{x m_1}{L_x} + \frac{y m_2}{L_y} \right)},
\]

and the discrete Fourier coefficients are

\[
\rho_n(m_1, m_2) = \frac{1}{\sqrt{(L_x - L_1)(L_y - L_2)}} \times \int_{L_1}^{L_1} \int_{L_2}^{L_2} e^{-2\pi i \left( \frac{x m_1}{L_x} + \frac{y m_2}{L_y} \right)} \zeta^n(x, y) d\xi d\eta.
\]

For any vector \( Z = (z_{1,1}, z_{1,2}, \ldots, z_{N_x,N_y})^T \), we define the discrete \( l^2 \) norm as follows

\[
||Z||_2^2 = \left( h_x h_y \sum_{i=1}^{N_x} \sum_{j=1}^{N_y} z_{jk}^2 \right)^{1/2}.
\]

So we have Parseval’s equality for the discrete Fourier transform, that is,

\[
||\zeta^n||_2^2 = h_x h_y \sum_{i=1}^{N_x} \sum_{j=1}^{N_y} |\zeta_{ij}^n|^2
= \int_{L_1}^{L_1} \int_{L_2}^{L_2} |\zeta^n(\xi, \eta)|^2 d\xi d\eta = \sum_{m_1=-\infty}^{\infty} \sum_{m_2=-\infty}^{\infty} |\rho_n(m_1, m_2)|^2.
\]
We can expand $\zeta_{jk}$ into Fourier series, and because the difference equations are linear, we can analyze the behavior of the total error by tracking the behavior of an arbitrary $n$th component. So we can suppose that the solution of Eq. (23) and Eq. (24) has the following form

$$\zeta_{jk}^n = \frac{1}{\sqrt{L_2 - L_1} \sqrt{L_4 - L_3}} \rho_n e^{i(\sigma_x L_4 + \sigma_y L_3)},$$

where $\sigma_x = 2m_1 \pi/(L_2 - L_1)$, $\sigma_y = 2m_2 \pi/(L_4 - L_3)$.

Substituting the above expression into Eq. (23) and Eq. (24) and for simplicity we let $s_1 = \sin^2 \frac{\sigma_x}{2}$, $s_2 = \sin^2 \frac{\sigma_y}{2}$, (thus $0 \leq s_i \leq 1$, $i = 1, 2$), then for $n = 1$ we have

$$(1 - 4 \left(\frac{1}{12} - \mu_x\right)s_1)(1 - 4 \left(\frac{1}{12} - \mu_y\right)s_2) \rho_1 = \left[16 \mu_x \mu_y s_1 s_2 + (1 - \frac{1}{3}s_1)(1 - \frac{1}{3}s_2)\right] \rho_0,$$

and for $2 \leq n \leq N$, we get

$$\begin{align*}
(1 - 4 \left(\frac{1}{12} - \mu_x\right)s_1)(1 - 4 \left(\frac{1}{12} - \mu_y\right)s_2) \rho_n &= \\
\left[32 \left(\frac{1}{144} + \mu_x \mu_y\right)s_1 s_2 + \frac{d_{n1} - d_{n1,2}}{d_{n1}} \left(- \frac{1}{3}(s_1 + s_2) + 1\right)\right] \rho_{n-1} \\
-16 \left(\frac{1}{144} + \mu_x \mu_y\right)s_1 s_2 + \frac{d_{n2} - d_{n2,1}}{d_{n2}} \left(- \frac{1}{3}(s_1 + s_2) + 1\right) \rho_{n-2} \\
+ \left(- \frac{1}{3}(s_1 + s_2) + 1\right) \left(\sum_{l=1}^{n-3} \left(\frac{d_{l+1} - d_{l1}}{d_{l1}}\right) \rho_l + \frac{d_{l1}}{d_{l1}} \rho_0\right).
\end{align*}$$

Consequently, we obtain

$$\rho_1 = \frac{16 \mu_x \mu_y s_1 s_2 + (1 - \frac{1}{3}s_1)(1 - \frac{1}{3}s_2)}{(1 - \frac{1}{3}s_1 + 4 \mu_x s_1)(1 - \frac{1}{3}s_2 + 4 \mu_y s_2)} \rho_0 \equiv \frac{\eta_0}{\mu} \rho_0,$$

$$\rho_n = \frac{32 \left(\frac{1}{144} + \mu_x \mu_y\right)s_1 s_2 + \frac{d_{n1} - d_{n1,1}}{d_{n1}} \left(- \frac{1}{3}(s_1 + s_2) + 1\right)}{(1 - \frac{1}{3}s_1 + 4 \mu_x s_1)(1 - \frac{1}{3}s_2 + 4 \mu_y s_2)} \rho_{n-1}$$

$$+ \frac{-16 \left(\frac{1}{144} + \mu_x \mu_y\right)s_1 s_2 + \frac{d_{n2} - d_{n2,2}}{d_{n2}} \left(- \frac{1}{3}(s_1 + s_2) + 1\right)}{(1 - \frac{1}{3}s_1 + 4 \mu_x s_1)(1 - \frac{1}{3}s_2 + 4 \mu_y s_2)} \rho_{n-2}$$

$$+ \left(1 - \frac{1}{3}(s_1 + s_2) + 1\right) \left(\sum_{l=1}^{n-3} \left(\frac{d_{l+1} - d_{l1}}{d_{l1}}\right) \rho_l + \frac{d_{l1}}{d_{l1}} \rho_0\right)$$

$$\equiv \sum_{l=1}^{n} \frac{\eta_{n-l}}{\mu} \rho_{n-l}, \quad 2 \leq n \leq N,$$

where $\mu = (1 - \frac{1}{3}s_1 + 4 \mu_x s_1)(1 - \frac{1}{3}s_2 + 4 \mu_y s_2)$. 

\[\text{Filomat 34:11 (2020), 3609–3626}\]
Proof. We will prove this claim by mathematical induction. For and the proof is completed.

Remark 3.1. The coefficient $\mu$ satisfies in the following relation $0 < \frac{1}{d} \leq \frac{9}{4}$. 

Now, we can give the following estimate.

Theorem 3.2. Suppose that $\rho_n (1 \leq n \leq N)$ are defined by (25) and (26), then for $\gamma \in (0, 1)$

$$|\rho_n| \leq \left(1 + \frac{C_T}{\mu}\right)^n |\rho_0|, \quad n = 1, 2, \cdots, N.$$  

Proof. We will prove this claim by mathematical induction. For $n = 1$ we have

$$|\rho_1| = \left| \frac{n_0}{\mu} \right| |\rho_0| \leq \left(1 + \frac{C_T}{\mu}\right) |\rho_0|.$$ 

Assume that

$$|\rho_m| \leq \left(1 + \frac{C_T}{\mu}\right)^m |\rho_0|, \quad m = 1, 2, \ldots, n - 1,$$

then we get

$$|\rho_n| = \left| \sum_{l=1}^{n} \frac{n_{n-l}}{\mu} \rho_{n-l} \right|$$

$$\leq \left| \frac{32}{1 + \frac{1}{d} + \mu s_1 \mu s_2 + \left( \frac{d_1 - d_{n-1}}{d_{n,n}} \right) - \left( \frac{1}{d} (s_1 + s_2) + 1 \right)} \right| \left| \left(1 + \frac{C_T}{\mu}\right)^{n-1} |\rho_0| \right|$$

$$+ \left| \frac{-16}{1 + \frac{1}{d} + \mu s_1 s_2 + \left( \frac{d_1 - d_{n-1}}{d_{n,n}} \right) - \left( \frac{1}{d} (s_1 + s_2) + 1 \right)} \right| \left| \left(1 + \frac{C_T}{\mu}\right)^{n-1} |\rho_0| \right|$$

$$+ \left| \frac{1 - \frac{1}{d} (s_1 + s_2)}{1 + \frac{1}{d} + \mu s_1 s_2 + \left( \frac{d_1 - d_{n-1}}{d_{n,n}} \right) - \left( \frac{1}{d} (s_1 + s_2) + 1 \right)} \right| \left| \left(1 + \frac{C_T}{\mu}\right)^{n-1} \sum_{l=1}^{n-3} \left( \frac{d_{l+1,n} - d_{l,n}}{d_{n,n}} \right) |\rho_0| + \frac{d_{l,n}}{d_{n,n}} |\rho_0| \right|$$

$$\leq \frac{48 \mu s_1 \mu s_2 + \left( \frac{1 - \frac{1}{d}}{1 + \frac{1}{d} + \mu s_1 s_2 + \left( \frac{d_1 - d_{n-1}}{d_{n,n}} \right) - \left( \frac{1}{d} (s_1 + s_2) + 1 \right)} \right)}{1 + \frac{1}{d} + \mu s_1 s_2 + \left( \frac{d_1 - d_{n-1}}{d_{n,n}} \right) - \left( \frac{1}{d} (s_1 + s_2) + 1 \right)} \left| \left(1 + \frac{C_T}{\mu}\right)^{n-1} |\rho_0| \right| \leq \left(1 + \frac{C_T}{\mu}\right)^n |\rho_0|.$$ 

and the proof is completed. \qed

Theorem 3.3. The compact finite difference scheme (21) is unconditionally stable for $\gamma \in (0, 1)$.

Proof. Suppose that $\hat{U}^n$ is the approximate solution of (21). Applying Theorem 3.2, Parseval’s equality and $n\tau \leq T$ we obtain
\[ \|U^n - \bar{U}^n\|_E^2 = \|\xi^n\|_E^2 \]
\[ = h_x h_y \sum_{k=1}^{N_x} \sum_{j=1}^{N_y} |e_{\nu j k}|^2 \]
\[ = \frac{h_x}{L_2 - L_1} \frac{h_y}{L_3 - L_4} \sum_{k=1}^{N_x} \sum_{j=1}^{N_y} |p_\nu e^{i \nu x + i \nu y + \nu \nu \mu I}|^2 \]
\[ = \frac{h_x}{L_2 - L_1} \frac{h_y}{L_3 - L_4} \sum_{k=1}^{N_x} \sum_{j=1}^{N_y} (1 + \frac{C_T}{\mu})^{2n} |p_0|^2 \]
\[ \leq \frac{h_x}{L_2 - L_1} \frac{h_y}{L_3 - L_4} \sum_{k=1}^{N_x} \sum_{j=1}^{N_y} (1 + \frac{C_T}{\mu})^{2n} \sum_{j=1}^{N_z} |p_0 e^{i \nu x + i \nu y + \nu \nu \mu I}|^2 \]
\[ = \left(1 + \frac{C_T}{\mu}\right)^{2n} \|\xi^n\|_E^2 \leq e^{2\mu C_T} \|U^0 - \bar{U}^0\|_E^2, \quad n = 1, 2, \ldots, N. \]

So that
\[ \|U^n - \bar{U}^n\|_E \leq e^{2\mu C_T} \|U^0 - \bar{U}^0\|_E, \]
which means that the scheme is unconditionally stable. \(\square\)

3.2. Convergence

We will use the Fourier analysis for the scheme (21).

For \(1 \leq j \leq N_x, 1 \leq k \leq N_y\), let \(e_{\nu j k}^0 = u_{\nu j k}^{\nu 0} - U_{\nu j k}^0\) and \(e_{\nu j k}^n = (e_{\nu j k}^{\nu 1}, e_{\nu j k}^{\nu 2}, \ldots, e_{\nu j k}^{\nu N_z})^T, R_{\nu j k}^n = (R_{\nu j k}^{\nu 1}, R_{\nu j k}^{\nu 2}, \ldots, R_{\nu j k}^{\nu N_z})^T, 1 \leq n \leq N. \) Then
\[ \left(1 + \left(\frac{1}{12} - \mu_x\right)\delta_x^2\right)\left(1 + \left(\frac{1}{12} - \mu_y\right)\delta_y^2\right) e_{\nu j k}^1 = R_{\nu j k}^1, \]
and for \(n \geq 2\)
\[ \left(1 + \left(\frac{1}{12} - \mu_x\right)\delta_x^2\right)\left(1 + \left(\frac{1}{12} - \mu_y\right)\delta_y^2\right) e_{\nu j k}^n = \left(\frac{1}{144} + \mu_x + \mu_y\right)(2e_{\nu j k}^{n-1} - e_{\nu j k}^{n-2}) \]
\[ + \frac{1}{d_{\nu j k}} \left(\delta_x^2 + \delta_y^2 + 1\right) \sum_{n=1}^{N_z} (d_{\nu j k} + d_{\nu j k}) e_{\nu j k}^n + R_{\nu j k}^n. \]

Similar to the stability we have
\[ \|e_{\nu j k}^n\|_E^2 = h_x h_y \sum_{\nu=1}^{N_x} \sum_{j=1}^{N_y} |e_{\nu j k}^n|^2 = \sum_{m_2=-\infty}^{\infty} \sum_{m_1=-\infty}^{\infty} |c_{\nu}(m_1, m_2)|^2, \]
and
\[ \|R_{\nu j k}^n\|_E^2 = h_x h_y \sum_{\nu=1}^{N_x} \sum_{j=1}^{N_y} |R_{\nu j k}^n|^2 = \sum_{m_2=-\infty}^{\infty} \sum_{m_1=-\infty}^{\infty} |d_{\nu}(m_1, m_2)|^2. \]
Now we assume that $c_{jk}^n$ and $R_{jk}^n$ are

$$c_{jk}^n = \frac{1}{\sqrt{L_2 - L_1} \sqrt{L_4 - L_3}} c_n e^{i(\sigma_z b_k + \sigma_\phi b_h)},$$

$$R_{jk}^n = \frac{1}{\sqrt{L_2 - L_1} \sqrt{L_4 - L_3}} d_n e^{i(\sigma_z b_k + \sigma_\phi b_h)},$$

where $\sigma_z = 2m_1 \pi / (L_2 - L_1)$ and $\sigma_\phi = 2m_2 \pi / (L_4 - L_3)$.

Substituting the above expression into Eqs. (27) and (28) and for simplicity we let $s_1 = \sin^2 \frac{\sigma_\phi h_x}{2}$, $s_2 = \sin^2 \frac{\sigma_\phi h_y}{2}$, (thus $0 \leq s_i \leq 1$, $i = 1, 2$), then for $n = 1$ we have

$$c_1 = \frac{1}{\mu} d_1,$$

and for $n \geq 2$

$$c_n = \frac{32 \left(\frac{1}{144} + \mu \frac{1}{144} \right) s_1 s_2 + \left(\frac{d_{n+1,n} - d_{n,n}}{d_{n,n}}\right) \left(- \frac{1}{2} (s_1 + s_2) + 1\right) c_{n-1}}{\mu} - 16 \left(\frac{1}{144} + \mu \frac{1}{144} \right) s_1 s_2 + \left(\frac{d_{n+1,n} - d_{n,n}}{d_{n,n}}\right) \left(- \frac{1}{2} (s_1 + s_2) + 1\right) c_{n-2}$$

$$+ \frac{1}{\mu} \left(\sum_{l=1}^{n-3} \left(\frac{d_{l+1,n} - d_{l,n}}{d_{n,n}}\right) c_l + \frac{d_{1,n}}{d_{n,n}} c_0 + \frac{d_n}{\mu}\right).$$

(32)

where $\mu = \left(1 - \frac{1}{2} s_1 + 4 \mu s_1\right) \left(1 - \frac{1}{2} s_2 + 4 \mu s_2\right)$.

Also from convergence of the series in the right hand side of (30), there is a positive constant $C_2$ such that

$$|d_n| \leq C_2 L \tau |d_1|, \quad n = 1, 2, \ldots, N.$$  (33)

**Theorem 3.4.** If $c_n$ be the solutions of equations (31) and (32), then there is positive constants $C_2$ and $B_n$ such that

$$|c_n| \leq C_2 B_n \left(1 + \frac{\tau L}{\mu}\right)^n |d_1|, \quad n = 1, 2, \ldots, N.$$  (34)

**Proof.** We use the mathematical induction for proof. For $n = 1$, from Eq. (31) we have

$$|c_1| \leq \frac{1}{\mu} |d_1| \leq \frac{\tau L}{\mu} C_2 |d_1| \leq C_2 \left(1 + \frac{\tau L}{\mu}\right) |d_1|.$$

Now, suppose that

$$|c_m| \leq C_2 B_m \left(1 + \frac{\tau L}{\mu}\right)^m |d_1|, \quad m = 1, 2, \ldots, n - 1.$$
Using Eq. (32), we have

\[ |c_n| \leq \left| \frac{32\left( \frac{1}{\mu^4} + \mu^4 \mu y \right) s_1 s_2 + \left( \frac{d_{n,1}-d_{n,2}}{d_{n,2}} \right) \left( -\frac{1}{3}(s_1 + s_2) + 1 \right) \right| C_2 B_{n-1} \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} |d_1| \]

\[ -16\left( \frac{1}{\mu^4} + \mu^4 \mu y \right) s_1 s_2 + \left( \frac{d_{n,1}-d_{n,2}}{d_{n,2}} \right) \left( -\frac{1}{3}(s_1 + s_2) + 1 \right) C_2 B_{n-2} \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} |d_1| \]

\[ + \frac{1 - \frac{1}{3}(s_1 + s_2)}{\mu} \sum_{l=1}^{n-3} \left( \frac{d_{l+1,n} - d_{l,n}}{d_{n,n}} \right) C_2 B_l \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} |d_1| \left| \frac{d_n}{\mu} \right|. \]

Now assume that

\[ C' = \max\{|B_1, B_2, \ldots, B_{n-1}|, \]

so

\[ |c_n| \leq \left| \frac{32\left( \frac{1}{\mu^4} + \mu^4 \mu y \right) s_1 s_2 + \left( \frac{d_{n,1}-d_{n,2}}{d_{n,2}} \right) \left( -\frac{1}{3}(s_1 + s_2) + 1 \right) \right| C_2 C' \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} |d_1| \]

\[ -16\left( \frac{1}{\mu^4} + \mu^4 \mu y \right) s_1 s_2 + \left( \frac{d_{n,1}-d_{n,2}}{d_{n,2}} \right) \left( -\frac{1}{3}(s_1 + s_2) + 1 \right) C_2 C' \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} |d_1| \]

\[ + \frac{1 - \frac{1}{3}(s_1 + s_2)}{\mu} \sum_{l=1}^{n-3} \left( \frac{d_{l+1,n} - d_{l,n}}{d_{n,n}} \right) C_2 C' \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} |d_1| + C_2 \left( 1 + \frac{\tau L}{\mu} \right)^n |d_1| \]

\[ \leq \frac{48\left( \frac{1}{\mu^4} + \mu^4 \mu y \right) s_1 s_2 + \left( -\frac{1}{3}(s_1 + s_2) + 1 \right) \left( 1 + \frac{\tau L}{\mu} \right)^{n-1} \left| d_1 \right| + \left( 1 + \frac{\tau L}{\mu} \right)^n |d_1| \]

\[ \leq C_2 C' \left( 1 + \frac{\tau L}{\mu} \right)^n |d_1| + C_2 \left( 1 + \frac{\tau L}{\mu} \right)^n |d_1| = C_2 (1 + C') \left( 1 + \frac{\tau L}{\mu} \right)^n |d_1| \]

\[ = C_2 B_n \left( 1 + \frac{\tau L}{\mu} \right)^n |d_1|. \]

This completes the proof. \( \square \)

**Theorem 3.5.** The compact ADI finite difference scheme (21) is convergent, and the order of convergence is \( O(\tau + h^4) \).

**Proof.** By Theorem (3.4) and using (29) and (30), we can obtain

\[ ||e^\tau||_e \leq C_2 B_n \left( 1 + \frac{\tau L}{\mu} \right)^n ||R^\tau||_e \leq C_2 B_n e^{\frac{\tau L}{\mu}} C_3 (\tau + h^4). \]

By remark (3.1) \( 0 \leq \frac{1}{\mu} \leq \frac{2}{3} \) and since \( n\tau \leq T \), we obtain

\[ ||e^\tau||_e \leq C (\tau + h^4), \quad C = C_2 B_n C_3 e^{\frac{\tau L}{\mu}}. \]

\( \square \)
4. Numerical Results

In this section, we perform some numerical calculations to test the convergence of our scheme. All the computations are performed using Matlab. In all examples we assume that \( h_x = h_y = h \) and use the error norm

\[
\|e^n\|_{\infty} = \max \left| u(x_j, y_k, t_n) - U^n_{jk} \right| : 0 \leq j \leq N_x, 0 \leq k \leq N_y,
\]

where \( e^n_{jk} = u^n_{jk} - U^n_{jk} \). Also we calculated the computational orders of the method presented in this paper with the following formula

\[
r(\tau, h) = \log_{2}(\|e(4\tau, 2h)\|_{\infty} / \|e(\tau, h)\|_{\infty}).
\]

**Example 4.1.** We first consider the following equation with the initial and boundary conditions

\[
\begin{align*}
\frac{\partial}{\partial t} u(x, y, t) &= \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(x, y, t), \quad 0 < x, y < 1, \ 0 \leq t \leq 1, \\
u(0, y, t) &= 0, u(1, y, t) = 0, \\
u(x, 0, t) &= 0, u(x, 1, t) = 0, \\
u(x, y, 0) &= 0.
\end{align*}
\]

We give exact solution \( u(x, y, t) = xy(1-x)(1-y)\sin(t) \), and for different \( \gamma \), we have different \( f(x, y, t) \). Numerical solution and pointwise errors have been demonstrated in Fig. (1). Table (1) gives the approximation errors and convergence rates for the fourth order compact difference scheme. We choose different space and time step sizes to obtain the numerical convergence rate in space.

![Numerical solution](image1)

![Pointwise errors](image2)

Figure 1: The plot of numerical solution and pointwise errors at \( \tau = 1/256 \) and \( h_x = h_y = 1/32 \) with \( \gamma = 0.2 \) for example 4.1.

**Example 4.2.** Consider the two dimensional time fractional diffusion equation as follows

\[
\begin{align*}
\frac{\partial^\gamma}{\partial t^\gamma} u(x, y, t) &= \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(x, y, t), \quad 0 < x, y < 1, \ 0 \leq t \leq 1, \\
u(0, y, t) &= \tau^{2}\gamma^{n}, u(1, y, t) = \tau^{2}\gamma^{n}, \\
u(x, 0, t) &= \tau^{2}\gamma^{n}, u(x, 1, t) = \tau^{2}\gamma^{n+1}, \\
u(x, y, 0) &= 0.
\end{align*}
\]

The exact solution for this problem is \( u(x, y, t) = \tau^{2}\gamma^{n+1} \). Table (2) gives the approximation errors and rates of \( u \) for the compact difference scheme of Example (4.2). Results show that the space rates are almost \( O(h^4) \) and are consistent with our theoretical analysis. Also Numerical solution and pointwise errors have been demonstrated in Fig. (2).
We give exact solution $u(3)$ in Table 1. The maximum absolute errors and their estimated convergence rates approximated by the compact ADI method are shown.

Consider the following equation

$$D_t^\gamma u(x, y, t) = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(x, y, t), \quad 0 < x, y < 1, \quad 0 \leq t \leq 1,$$

$$u(0, y, t) = 0, \quad u(1, y, t) = 0,$$

$$u(x, 0, t) = 0, \quad u(x, 1, t) = 0,$$

$$u(x, y, 0) = \sin(\pi x) \sin(\pi y).$$

We give exact solution $u(x, y, t) = (t^2 + 1) \sin(\pi x) \sin(\pi y)$, and for different $\gamma$, we have different $f(x, y, t)$. The maximum absolute errors and their estimated convergence rates approximated by the compact ADI method are shown in Table 1. The numerical results of Example 4.3 are provided to show that the proposed approximation method is computationally efficient.

### Table 1: Error and experiment order of convergence for different values of $\gamma$, for example 4.1

| $\gamma$ | $h$ | $\tau$ | $\|e^n\|_{\infty}$ | Order |
|----------|-----|--------|------------------|-------|
| 0.1      | 1/3 | 1/3    | 1.3683e - 02     | -     |
| 0.1      | 1/4 | 1/4    | 7.0627e - 04     | 4.2760|
| 0.1      | 1/5 | 1/5    | 5.6257e - 05     | 3.6501|
| 0.1      | 1/6 | 1/6    | 3.3538e - 06     | 4.0682|
| 0.3      | 1/3 | 1/3    | 9.5564e - 03     | -     |
| 0.3      | 1/4 | 1/4    | 2.8612e - 04     | 5.0618|
| 0.3      | 1/5 | 1/5    | 3.9921e - 05     | 2.4989|
| 0.3      | 1/6 | 1/6    | 2.3441e - 06     | 4.0900|
| 0.5      | 1/3 | 1/3    | 5.0387e - 03     | -     |
| 0.5      | 1/4 | 1/4    | 4.6054e - 04     | 3.4516|
| 0.5      | 1/5 | 1/5    | 2.4932e - 05     | 4.2072|
| 0.5      | 1/6 | 1/6    | 1.6940e - 06     | 3.8795|
| 0.7      | 1/3 | 1/3    | 1.0617e - 03     | -     |
| 0.7      | 1/4 | 1/4    | 3.4366e - 04     | 1.6273|
| 0.7      | 1/5 | 1/5    | 1.7136e - 05     | 4.3258|
| 0.7      | 1/6 | 1/6    | 1.0063e - 06     | 4.0899|
| 0.9      | 1/3 | 1/3    | 1.2361e - 03     | -     |
| 0.9      | 1/4 | 1/4    | 1.0447e - 04     | 3.5646|
| 0.9      | 1/5 | 1/5    | 1.6022e - 05     | 4.2301|
| 0.9      | 1/6 | 1/6    | 3.3166e - 07     | 5.5942|

Figure 2: The plot of numerical solution and pointwise errors at $\tau = 1/256$ and $h_x = h_y = 1/32$ with $\gamma = 0.2$ for example 4.2.

### Example 4.3

Consider the following equation

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(x, y, t), \quad 0 < x, y < 1, \quad 0 \leq t \leq 1,$$

$$u(0, y, t) = 0, \quad u(1, y, t) = 0,$$

$$u(x, 0, t) = 0, \quad u(x, 1, t) = 0,$$

$$u(x, y, 0) = \sin(\pi x) \sin(\pi y).$$
In this example the new compact ADI scheme is used to solve the following equation

Example 4.4. In this example the new compact ADI scheme is used to solve the following equation

\[
\begin{align*}
\text{CF} \frac{\partial}{\partial \tau} u(x, y, t) &= \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(x, y, t), \quad 0 < x, y < 1, \quad 0 \leq t \leq 1, \\
u(0, y, t) &= 0, u(1, y, t) = 0, \\
u(x, 0, t) &= 0, u(x, 1, t) = 0, \\
u(x, y, 0) &= x^2(1 - x)^3 \sin(\pi y)e^{\mu}\gamma.
\end{align*}
\]

The exact solution for this problem is \( u(x, y, t) = (t^2 + 1)x^2(1 - x)^3 \sin(\pi y) e^{\mu+2\gamma} \). Numerical results for time and space with \( \gamma = 0.1, 0.3, 0.5, 0.7, 0.9 \) are presented in Table (4), respectively. Figure (4) shows the numerical solutions and pointwise errors.

![Figure 3: The plot of numerical solution and pointwise errors at \( \tau = 1/256 \) and \( h_x = h_y = 1/32 \) with \( \gamma = 0.2 \) for example 4.3.](image)

| \( \gamma \) | \( h \) | \( \tau \) | \( ||e||_{\infty} \) | Order |
|---|---|---|---|---|
| 0.1 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 2.2886e - 02 | - |
| 0.1 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.2102e - 03 | 4.2411 |
| 0.1 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 8.9767e - 05 | 3.7529 |
| 0.1 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 5.5718e - 06 | 4.0010 |
| 0.3 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.9323e - 02 | - |
| 0.3 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.0779e - 03 | 4.1640 |
| 0.3 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 6.6135e - 05 | 4.0267 |
| 0.3 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 4.1376e - 06 | 3.9985 |
| 0.5 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.4057e - 02 | - |
| 0.5 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 7.0760e - 04 | 4.3122 |
| 0.5 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 4.0324e - 05 | 4.1332 |
| 0.5 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 2.5167e - 06 | 4.0020 |
| 0.7 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 5.2053e - 03 | - |
| 0.7 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.2603e - 04 | 5.3681 |
| 0.7 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 6.2757e - 06 | 4.3278 |
| 0.7 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 3.6523e - 07 | 4.1029 |
| 0.9 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.6284e - 02 | - |
| 0.9 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 1.551e - 03 | 3.8174 |
| 0.9 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 7.6274e - 05 | 3.9207 |
| 0.9 | \(-\frac{1}{10} \) | \(-\frac{1}{10} \) | 4.8040e - 06 | 3.9889 |
Table 3: Error and experiment order of convergence for different values of $\gamma$, for example 4.3

| $\gamma$ | $h$ | $\tau$ | $\|e\|_{\infty}$ | Order |
|----------|-----|--------|-------------------|-------|
| 0.1      | 1/4 | 1/2    | 5.3409e-01        | -     |
| 0.1      | 1/8 | 1/4    | 2.8900e-02        | 4.279 |
| 0.1      | 1/16| 1/8    | 2.1592e-03        | 3.742 |
| 0.1      | 1/32| 1/16   | 1.5036e-04        | 3.844 |
| 0.3      | 1/4 | 1/2    | 4.7311e-01        | -     |
| 0.3      | 1/8 | 1/4    | 2.7209e-02        | 4.120 |
| 0.3      | 1/16| 1/8    | 1.9246e-03        | 3.821 |
| 0.3      | 1/32| 1/16   | 1.7690e-04        | 3.443 |
| 0.5      | 1/4 | 1/2    | 3.9814e-01        | -     |
| 0.5      | 1/8 | 1/4    | 2.2690e-02        | 4.133 |
| 0.5      | 1/16| 1/8    | 1.7522e-03        | 3.694 |
| 0.5      | 1/32| 1/16   | 2.1784e-04        | 3.007 |
| 0.7      | 1/4 | 1/2    | 3.0489e-01        | -     |
| 0.7      | 1/8 | 1/4    | 1.5001e-02        | 4.345 |
| 0.7      | 1/16| 1/8    | 1.4322e-03        | 3.388 |
| 0.7      | 1/32| 1/16   | 2.2312e-04        | 2.682 |
| 0.9      | 1/4 | 1/2    | 1.8515e-01        | -     |
| 0.9      | 1/8 | 1/4    | 5.1344e-03        | 5.172 |
| 0.9      | 1/16| 1/8    | 3.5669e-04        | 3.847 |
| 0.9      | 1/32| 1/16   | 4.4063e-05        | 3.017 |

Figure 4: The plot of numerical solution and pointwise errors at $\tau = 1/256$ and $h_x = h_y = 1/32$ with $\gamma = 0.2$ for example 4.4.

5. Conclusions

In this paper, we constructed a high-order compact alternating direction implicit method for the solution of two dimensional time-fractional diffusion equation with Caputo-Fabrizio derivative. The time fractional derivative of the mentioned equation approximated by a scheme of order $O(\tau^2)$ and spatial derivatives replaced with a fourth order compact difference scheme. We prove that the scheme is unconditionally stable for $\gamma \in (0, 1)$. Numerical results confirmed the theoretical results of the proposed scheme, i.e the scheme has fourth order of accuracy in space.
Table 4: Error and experiment order of convergence for different values of $\gamma$, for example 4.4

| $\gamma$ | $h$ | $\tau$ | $\|e^n\|_{\infty}$ | Order |
|----------|-----|--------|-------------------|-------|
| 0.1      | 9.7182e - 02 | 3.6109 |
| 0.1      | 7.9542e - 03 | 3.6109 |
| 0.1      | 5.3281e - 04 | 3.6109 |
| 0.1      | 3.6008e - 05 | 3.6109 |
| 0.3      | 8.7864e - 02 | 3.6109 |
| 0.3      | 6.4657e - 03 | 3.6109 |
| 0.3      | 4.4755e - 04 | 3.6109 |
| 0.3      | 3.4778e - 05 | 3.6109 |
| 0.5      | 75959e - 02 | 3.6109 |
| 0.5      | 5.0152e - 03 | 3.6109 |
| 0.5      | 3.6417e - 04 | 3.6109 |
| 0.5      | 3.6070e - 05 | 3.6109 |
| 0.7      | 6.0359e - 02 | 3.6109 |
| 0.7      | 3.1687e - 03 | 3.6109 |
| 0.7      | 2.6388e - 04 | 3.6109 |
| 0.7      | 3.3105e - 05 | 3.6109 |
| 0.9      | 4.0009e - 02 | 3.6109 |
| 0.9      | 1.1959e - 03 | 3.6109 |
| 0.9      | 7.1557e - 05 | 3.6109 |
| 0.9      | 6.8432e - 06 | 3.6109 |

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