EXTRAGALACTIC VERY HIGH ENERGY GAMMA-RAY BACKGROUND

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ABSTRACT

We study the origin of the extragalactic diffuse gamma-ray background using the data from the Fermi telescope. To estimate the background level, we count photons at high Galactic latitudes $|b| > 60^\circ$. Subtracting photons associated with known sources and the residual cosmic-ray and Galactic diffuse backgrounds, we estimate the extragalactic gamma-ray background (EGB) flux. We find that the spectrum of EGB in the very high energy band above 30 GeV follows the stacked spectrum of BL Lac objects. Large Area Telescope data reveal the positive $(1 + z)^k$, $1 < k < 4$ cosmological evolution of the BL Lac source population consistent with that of their parent population, Fanaroff–Riley type I radio galaxies. We show that EGB at $E > 30$ GeV could be completely explained by emission from unresolved BL Lac objects if $k \simeq 3$.

Key words: BL Lacertae objects: general – gamma rays: diffuse background

Online-only material: color figures

1. INTRODUCTION

The study of diffuse background emission produced by faint sources with flux levels below the sensitivity of a telescope is commonly used to constrain the nature of source populations in the universe and their cosmological evolution. In the high-energy (HE) $\gamma$-ray band (0.1–100 GeV) the diffuse extragalactic $\gamma$-ray background (EGB) was detected for the first time by the SAS-2 satellite (Fichtel et al. 1978), was further studied by EGRET telescope on board the CGRO mission (Sreekumar et al. 1998; Strong et al. 2004), and, most recently, by the Large Area Telescope (LAT) on board the Fermi satellite (Abdo et al. 2010a).

It is often assumed that the dominant contribution to the EGB is given by distant active galactic nuclei (AGNs), in particular, blazars (Padovani et al. 1993; Stecker et al. 1993; Chiang et al. 1995; Stecker & Salamon 1996; Mukherjee & Chiang 1999; Mücke & Pohl 2000; Inoue & Totani 2009). However, a recent study by the Fermi Collaboration reveals that blazars might contribute only a relatively small fraction of the HE EGB level (Abdo et al. 2010b), while a significant part of the EGB should be either explained by a yet unknown source population or have a truly diffuse nature (see, however, Stecker & Venters 2011).

The EGB in the very high energy (VHE; $\gamma$-rays in the $E \gtrsim 100$ GeV) range has never been measured. On one hand, the effective collection area of previous space-based $\gamma$-ray telescopes was not sufficient to achieve significant photon statistics in this energy band. On the other hand, the efficiency of cosmic-ray background rejection in the ground-based Cherenkov telescopes, such as HESS, MAGIC, and VERITAS, is not sufficient for the detection of the isotropic diffuse EGB on top of the cosmic-ray background. Thus, the properties and the origin of VHE EGB remain largely unconstrained until now.

It is clear that the VHE EGB should contain a contribution from the unresolved point sources. The main candidate source class is, as in the case of HE EGB, that of blazars. At the same time, the VHE EGB could contain, apart from the contribution from unresolved extragalactic point sources, genuine diffuse components which could be produced via several mechanisms. For example, if the spectra of a large number of $\gamma$-ray-loud AGNs extend to the energies above 300 GeV, all of the power initially emitted in $\gamma$-rays with energies higher than $\sim 300$ GeV is absorbed in the pair production of $\gamma$-rays on the cosmological infrared and/or microwave backgrounds (Gould & Schreder 1967). Secondary inverse Compton (IC) emission of electron–positron pairs deposited in the intergalactic space as a result of the pair production leads to the generation of diffuse extragalactic emission in the VHE energy band (Coppi & Aharonian 1997). Another mechanism which can lead to the generation of diffuse component of VHE EGB is electromagnetic cascade initiated in the intergalactic space by ultrahigh energy cosmic rays interacting with cosmic microwave background photons (Berezinsky & Smirnov 1975; Kalashev et al. 2004; Berezinsky et al. 2011). The cascade channels the power from the highest energies of about $10^{20}$ eV down to the ~0.1 GeV band in which the mean free path of the $\gamma$-rays becomes comparable to the size of the visible part of the universe.

Apart from the “guaranteed” (but, possibly, very weak) diffuse contributions, isotropic VHE $\gamma$-ray background might contain contributions from “exotic” diffuse sources, like diffuse emission from annihilation of Dark Matter particles in the outer halo of the Milky Way galaxy and the annihilation signal accumulated from the dark matter halos of all galaxies in the course of cosmological evolution (Abdo et al. 2010c).

Whatever the sources of VHE EGB are, they are scattered across the universe, so that a significant contribution to the flux is produced at redshifts $z \sim 1$. A known effect of the absorption of VHE $\gamma$-rays due to the interactions with infrared/optical extragalactic background light (EBL) should lead to attenuation of the $E > 50$ GeV signal produced by the sources at large redshifts $z \sim 1$ (Gould & Schreder 1967; Kneiske et al. 2004; Franceschini et al. 2008; Stecker et al. 2006; Gilmore et al. 2009). This should leave an “imprint” on the VHE EBL spectrum, which should have the form of a gradual suppression with the increasing photon energy. Detecting an
EBL suppression feature in the EGB spectrum would provide an important constraint on the (largely uncertain) evolution of the EBL density and spectrum up to redshifts $z \sim 1$. Such a constraint is otherwise difficult to obtain from the studies of individual extragalactic VHE $\gamma$-ray sources because of the limited signal statistics at the highest energies, especially for the sources at significant redshifts.

In the following, we discuss the measurement of the EGB derived from the data of Fermi/LAT telescope (Atwood et al. 2009). The measurement is obtained from the counting of photons at high Galactic latitudes, after the subtraction of the Galactic diffuse emission and the residual cosmic-ray background not rejected by the LAT data analysis software. We compare the measurement of EGB obtained in this way with the measurement previously derived from the likelihood analysis of all-sky data by Abdo et al. (2010a).

The EGB flux above 30 GeV is comparable to the flux in extragalactic VHE $\gamma$-ray sources resolved by Fermi. We find that the spectrum of EGB in this energy range follows the cumulative spectrum of the resolved sources. Dominant population of extragalactic VHE $\gamma$-ray sources is BL Lac objects. Noticing the similarity of the spectrum of VHE EGB and of the cumulative BL Lac VHE $\gamma$-ray spectrum, we put forward a hypothesis that the VHE EBL is produced by unresolved BL Lac objects with fluxes below the sensitivity of LAT. We explore this hypothesis and show that it could be valid if BL Lac objects follow a positive cosmological evolution pattern, characteristic for other types of AGNs, in particular for the parent population of BL Lac objects, Fanaroff–Riley type I (FR I) radio galaxies.

2. DATA SELECTION AND DATA ANALYSIS

For our analysis, we consider all publicly available LAT data from 2008 August 4 to 2011 January 23. We process the data using Fermi Science Tools. We filter the entire data set with gtselect and gtmktime tools following the recommendations of Fermi team and retain only events belonging to “ultraclean” (P7ULTRACLEAN_V6) event class, which has minimal residual cosmic-ray contamination.

In order to estimate of the contribution of point sources to the total flux requires the separation of the photons coming from the point sources from those produced by the diffuse emission. Such separation is most straightforward for the photons with narrow point-spread function (PSF). Taking this into account, we select two sub-classes (which provide dominant contribution above 100 GeV.) Extrapolation of the IC component gives a sub-dominant contribution to the ultraclean events (which provide dominant contribution to the ultraclean events) with the most compact PSF, the sub-classes selected by imposing the selection criterion EVENT\_CLASS=65311 or 32543. Other sub-classes of the ultraclean events have worse PSF. Point-source contribution in these photons suffers from an additional uncertainty. Taking this into account, we restrict our attention to the subset of the ultraclean events with the best PSF.

We retain events with an Earth zenith angle $\theta_e \leq 100^\circ$. To estimate the flux from the photon counts we use gtxposure tool. We consider only events at high Galactic latitudes in the regions $|b| \geq 60^\circ$.

Our analysis is based on the so-called Pass 7 selection of the LAT data (see http://fermi.gsfc.nasa.gov/ssc/data/access/). However, we use a comparison of the Pass 7 data with the previous Pass 6 data selection in the estimate of residual cosmic-ray contamination of the set of events chosen for the analysis.

The residual cosmic-ray fraction in the Pass 6 data was studied in detail by Abdo et al. (2010a). It is possible to re-calculate the residual cosmic-ray fraction for any new selection of events, including the one considered in our analysis, in a straightforward way, as explained in Section 3.3.

3. DIFFUSE $\gamma$-RAY BACKGROUND

The signal detected by LAT at high Galactic latitudes contains four types of contributions: emission from point sources, diffuse $\gamma$-ray emission from the Galaxy, EGB, and residual cosmic-ray background not rejected by the analysis software. To measure the EGB flux, one must separate the contributions from the four components in the overall signal in a given energy band.

3.1. Point-source Contribution

The point-source component could be singled out in a straightforward way if the set of sources detectable in a given energy band is known. To define the set of sources, we find the sources correlating with the directions of the arrival of photons in each energy band, using the method described by Neronov et al. (2011). To calculate the total number of photons associated with the sources, we construct a cumulative distribution of photons as a function of the distance $\theta$ from the source and split it on the background and source contributions. The background contribution grows asymptotically as $\theta^2$, while the source contribution asymptotically reaches constant. An example of the cumulative photon distribution around the source positions in the 12.5–25 GeV energy band is shown in Figure 1.

3.2. Galactic Diffuse Emission Contribution

The contribution of the diffuse emission from the Galaxy should be found via a detailed fitting of the all-sky photon distribution to an all-sky spatial and spectral template. This contribution is best constrained by the all-sky photon distribution in the 0.1–10 GeV energy band, where event statistics are very high. Detailed fitting of the Galactic diffuse emission to the data in the 0.1–10 GeV band was done by Abdo et al. (2010a). In our analysis, we rely on the best-fit model of Galactic diffuse emission derived by Abdo et al. (2010a). This model is available in the sky region of interest, $|b| \geq 60^\circ$ (see Figure 6 of the Supplemental Material in Abdo et al. 2010a). The uncertainties of the Galactic diffuse emission model are also discussed by Abdo et al. (2010a). We take these uncertainties into account.

The model consists of two main contributions: the “atomic hydrogen” component produced by interactions of cosmic rays with interstellar matter and “IC” component produced by IC emission from cosmic-ray electrons. Extrapolation of the atomic hydrogen component to the highest LAT energies is straightforward: the pion decay spectrum follows the cosmic-ray spectrum and extrapolation has the form of a simple power law with photon index $\sim 2.7$, the same as the slope of the cosmic-ray spectrum. This component gives a sub-dominant contribution above 100 GeV. Extrapolation of the IC component depends on the unknown shape of the (average over interstellar medium) cosmic-ray electron spectrum at the energies above TeV. We have checked that in the model of Abdo et al. (2010a) the spectrum of IC component is consistent with the
Figure 1. Cumulative front photon distributions around point sources in 12.5–25 GeV energy band. The red points show the source photons and blue points show the background. The horizontal lines (from top to bottom) show the 100%, 95%, and 68% levels.
(A color version of this figure is available in the online journal.)

Figure 2. Extrapolation of the spectrum of the Galactic diffuse emission at high Galactic latitudes $|b| \geq 60^\circ$ to the 100 GeV energy range. The blue and cyan dotted lines below 100 GeV show the contributions from cosmic-ray interactions with interstellar medium and inverse Compton scattering from cosmic electrons calculated by Abdo et al. (2010a). The continuation of the cyan dotted line above 100 GeV is calculated assuming that inverse Compton emission is produced by electrons with a cutoff power-law spectrum with cutoff at 1 TeV. The gray dashed line is the sum of the cosmic-ray and inverse Compton contributions. The solid gray line shows the overall diffuse emission spectrum in which the inverse Compton emission is produced by electron distribution without high-energy cutoff at 1 TeV. The red thick solid line shows the spectrum used for subtraction of Galactic component from the overall high Galactic latitude diffuse emission flux.
(A color version of this figure is available in the online journal.)

The HE cutoff of the local cosmic-ray electron spectrum is most likely determined by the distance to the closest cosmic-ray electron sources, e.g., to the closest pulsar wind nebulae (Aharonian 2004), rather than by the intrinsic cutoff in the

spectrum of IC scattering of the local interstellar radiation field (Moskalenko et al. 2006) by electrons with the spectrum $dN_e/dE \sim E^{-3} \exp(-E/1\text{ TeV})$. This electron spectrum is consistent with the cosmic-ray electron spectrum observed on Earth (Ackermann et al. 2010; Aharonian et al. 2008). The IC spectrum produced by such electron population is shown by the cyan dotted line in Figure 2. The overall Galactic diffuse emission spectrum at high Galactic latitudes is then the sum of the atomic hydrogen and IC contributions, shown by the dashed black line in Figure 2.
injection spectrum of electrons from the sources. This means that the local measurement of the HE cutoff of the cosmic electron spectrum does not provide a measurement of the HE cutoff in the injection spectrum of electrons. It is possible that Galactic cosmic electron sources inject electrons with energies much higher than \( \sim 1 \) TeV. This possibility is shown by the solid gray line in Figure 2 which shows the sum of the atomic hydrogen contribution with the IC emission from electrons without HE cutoff in the spectrum. The IC component still exhibits suppression at the energies \( \sim 1 \) TeV because of the Klein–Nishina effect.

To take into account the above-mentioned uncertainty of the IC component, we adopt an approximation \( dN_\gamma/dE \sim E^{-2.5} \exp(-E/2 \text{ TeV}) \) for the high Galactic latitude diffuse emission spectrum at \( E > 100 \) GeV. This approximation is shown by the red thick solid line in Figure 2. This spectrum lies exactly in the middle of the two extreme possibilities: TeV-scale HE cutoff in the cosmic electron spectrum and no cut-off in the cosmic electron spectrum. One should take into account that the uncertainty of this approximation reaches \( \sim 50\% \) at the highest energies. We take this uncertainty into account in the calculation of the EGB spectrum by adding it as a systematic error. The two extreme possibilities for the behavior of electron spectrum above 1 TeV (exponential cut-off exactly at 1 TeV and no cut-off at all) provide a good estimate of the overall uncertainty of electron spectrum in the interstellar medium in this range. The uncertainty of the shape of electron spectrum dominates the uncertainty of the IC component of Galactic diffuse emission at high Galactic latitudes.

### 3.3. Residual Cosmic-ray Background Contribution

To estimate the residual cosmic-ray background in the set of events selected for the analysis, we rely on the knowledge of residual. The residual cosmic-ray background in the dataclean event class of Pass 6 data is extensively discussed in Abdo et al. (2010a). The residual cosmic-ray background in the subset of Pass 7 superclean events used in our analysis could be calculated from the known residual cosmic-ray background in the Pass 6 dataclean events via a straightforward comparison of statistics of events on- and off-point sources in the two classes.

First, the residual cosmic-ray fraction in the Pass 6 dataclean events should be calculated from the known suppression factor of cosmic-ray event at transition from the diffuse event class to the dataclean event class in Pass 6 (see Abdo et al. 2010a for the detailed discussion of the suppression factor). In each of the two event classes, the entire event set consists of a certain number of \( \gamma \)-ray events \( N_{\gamma,i} \) and a certain number of residual cosmic-ray events, \( N_{\text{CR},i} \), where \( i \) stands for 3 + 4 or 4. Cleaning of the event set done to produce the dataclean event set from diffuse set results in the rejection of a large fraction of the cosmic-ray events, \( \alpha_{\text{CR}} \), with \( \alpha_{\text{CR}} \ll 1 \). However, it results also in a rejection of a number of true \( \gamma \)-ray events, so that \( N_{\gamma,i} = \alpha_{\gamma} N_{\gamma,i+4} \) with \( \alpha_{\gamma} < 1 \).

The suppression factor \( \alpha_{\text{CR}} \) is known as a function of energy from the Monte Carlo simulations of cosmic ray and \( \gamma \)-ray induced events in the LAT detector by Abdo et al. (2010a). The suppression factor \( \alpha_{\gamma} \) could be found directly from the data set, by comparing statistics of events coming from the point sources in the diffuse and dataclean event classes (see Section 3.1). In the calculation of \( \alpha_{\gamma} \) all the photons associated with \( \sim 10^3 \) point sources listed in the Fermi two-year catalog (Abdo et al. 2010d) could be used. This provides very large event statistics so that uncertainty of \( \alpha_{\gamma} \) is negligible. Knowing the total numbers of events in the two event classes \( N_{\text{tot},i} \) one can resolve the system of equations

\[
\begin{align*}
N_{\text{CR},4}/\alpha_{\text{CR}} + N_{\gamma,4}/\alpha_{\gamma} &= N_{\text{tot},3+4} \\
N_{\text{CR},4} + N_{\gamma,4} &= N_{\text{tot},4} 
\end{align*}
\]

(1)

with respect to \( N_{\gamma,4}, N_{\text{CR},4} \) to find the residual cosmic-ray background in each energy bin for the dataclean event class.

The residual cosmic-ray fraction in the sub-class of the Pass 7 superclean events used in our analysis is then estimated in a similar way, once the residual cosmic-ray fraction \( \kappa_4 \) in the point-source-subtracted set of the Pass 6 events, \( N_{\text{CR,off},4} = \kappa_4 N_{\text{off},4} \), is known.

Indeed, the transition from the Pass 6 dataclean events to the Pass 7 events belonging to the event classes 65311 and 32543 leaves a fraction \( \alpha_{\gamma,6\rightarrow7} \) of \( \gamma \)-ray events (actually, \( \alpha_{\gamma,6\rightarrow7} > 1 \) in a broad energy range around 10 GeV). It also suppresses or increases the residual cosmic-ray background, so that the residual cosmic-ray fraction in the off-source events changes from \( \kappa_4 \) to \( \kappa_7 \). The off-source events in the two classes are then the sum of the diffuse \( \gamma \)-ray emission photons and of the residual cosmic rays:

\[
\begin{align*}
\kappa_4 N_{\text{off},4} + N_{\gamma,\text{off},4} &= N_{\text{off},4} \\
\kappa_7 N_{\text{off},7} + \alpha_{\gamma,6\rightarrow7} N_{\gamma,\text{off},4} &= N_{\text{off},7}.
\end{align*}
\]

(2)

Knowing the statistics of the off-source events in the Pass 6 and Pass 7 events, \( N_{\text{off},4} \) and \( N_{\text{off},7} \), one could find the residual cosmic-ray fraction in the Pass 7 data

\[
\kappa_7 = 1 - \alpha_{\gamma,6\rightarrow7}(1 - \kappa_6) \frac{N_{\text{off},4}}{N_{\text{off},7}}.
\]

(3)

The resulting estimates of the level of residual cosmic-ray background for the events selected in the Pass 7 data in energy bins between 3 and 100 GeV are shown by the gray data points in Figure 3.

From Figure 3 one could see that the contribution of the residual cosmic rays to the signal at 100 GeV is likely to be small. However, extrapolation of the estimate of efficiency of rejection of the residual cosmic-ray background much above 100 GeV is highly uncertain. It is possible that the efficiency of rejection of both the nuclear and electron/positron component of the cosmic-ray flux drops because of the similarity of the cosmic ray and \( e^+e^- \) pair tracks with large Lorentz factors. Inefficient rejection of the residual cosmic rays might lead to the contamination of the diffuse background signal and lead to a large overestimation of the diffuse background flux. Because of this problem, we are able to only derive an upper limit on the EGB at the energies much above 100 GeV (for the energy band at 100 GeV we show a comparison between the 95% confidence level upper limit and the measurement). A proper measurement of the EGB flux at the highest energies accessible to LAT would require extensive Monte Carlo simulations taking into account detector response (M. Ackermann 2010, private communication).

### 3.4. Extragalactic \( \gamma \)-ray Background Spectrum

EGB flux could be found by subtracting the point source, Galactic diffuse, and residual cosmic-ray contributions to the
Figure 3. Estimate of the flux of isotropic component of diffuse emission obtained by the direct photon counting method (red thick line, data points, and upper limits). For comparison, the spectrum of isotropic component of diffuse sky emission, obtained using likelihood analysis at lower energies by Abdo et al. (2010a), is shown as a red shaded region. The pink upper limits above 100 GeV are from M. Ackermann (2010, private communication). Black data points show the total point-source-subtracted flux from the North and South Galactic Pole regions at $|b| \geq 60^\circ$. The solid line error bars show statistical error. The dashed line error bars show the systematic error at the level of $\approx 20\%$ stemming from the uncertainty of the Instrument Response Functions (IRFs) of LAT (see http://fermi.gsfc.nasa.gov/ssc/data/analysis/LAT_caveats.html). The black horizontally shaded region shows the point-source-subtracted flux in the Galactic Pole regions found by Abdo et al. (2010a). The gray data points and gray curve show the estimate of the residual cosmic-ray background in the event set used in this analysis. The Galactic diffuse emission contribution to the total flux at high Galactic latitude could not be negligibly small in the 100 GeV band. Taking this into account, it is not surprising that our estimate of the VHE EGB flux is somewhat lower than the total diffuse emission flux at high Galactic latitudes and is, respectively, lower than the upper limit derived by M. Ackermann (2010, private communication). It is useful to note that extrapolating the EGB spectrum as a power-law spectrum to $E \geq 100$ GeV band would give the spectrum consistent with the data above 100 GeV. With the current LAT exposure, there is still no evidence for suppression of VHE EGB flux due to absorption on EBL. A larger exposure time is needed to verify the presence of the feature. As it is mentioned in the Introduction, suppression of the flux above 100 GeV due to the absorption of VHE $\gamma$-rays on the EBL is expected if the EGB is accumulated over the cosmological distance scale. Detection of such suppression would be an important test of the origin of EGB.

As a matter of fact, the level of Galactic diffuse emission at high Galactic latitudes is comparable to the level of EGB in the entire energy range $E > 10$ GeV. The Galactic diffuse emission contribution to the total flux at high Galactic latitude could not be negligibly small in the 100 GeV band. Taking this into account, it is not surprising that our estimate of the VHE EGB flux is somewhat lower than the total diffuse emission flux at high Galactic latitudes and is, respectively, lower than the upper limit derived by M. Ackermann (2010, private communication).

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There is no a priori reason why the two fluxes should be nearly equal. Thus, the equality of the two contributions poses a “fine-tuning” problem which requires further investigation.
4. \( E > 30 \text{ GeV} \) EXTRAGALACTIC \( \gamma \)-RAY BACKGROUND FROM POINT SOURCES

As was mentioned in the Introduction, different types of point and diffuse sources could contribute to the EGB in the VHE band. The main class of extragalactic point sources detected by \textit{Fermi} is blazars, which are divided into two subclasses: BL Lac type objects and Flat Spectrum Radio Quasars (FSRQs). Over the first year of operation, the LAT has detected some \( \sim 700 \) such objects above 100 MeV energy (Abdo et al. 2010b). BL Lac objects and FSRQs have somewhat different spectral characteristics in the \( \gamma \)-ray band, with the spectra of BL Lac objects being systematically harder than the spectra of FSRQs (Abdo et al. 2010b). The hardness of the spectra of BL Lac objects implies that they might produce a significant contribution to the overall \( \gamma \)-ray flux in the VHE band. In fact, most of the extragalactic VHE \( \gamma \)-ray sources detected until now by the ground-based \( \gamma \)-ray telescopes sensitive above 100 GeV are BL Lac objects.9

Figure 4 shows the breakdown of the point-source contributions to the high Galactic latitude flux by source type. One could clearly see that the dominant contribution is given by BL Lac objects that provide \( \geq 90\% \) of the total point-source flux above 30 GeV. The cumulative spectrum of the other major blazar class, FSRQ, has an HE cutoff at \( \sim 10 \text{ GeV} \) so that FSRQ contribution to the point-source flux is negligible in the VHE band. From Figure 4, one could see that the total point-source flux calculated from the cumulative photon distribution around stacked point sources in the high Galactic latitude regions (see Section 3.1) is in a good agreement with the total point-source flux calculated using the likelihood analysis by Abdo et al. (2010a), as shown by the green shaded region. Above 50 GeV, 90\% of source photons come from BL Lac objects and 10\% from “Other” \textit{Fermi} sources, which are dominated by unidentified sources with some contribution from nearby AGNs.

The EGB spectral shape above 30 GeV follows the cumulative point-source spectrum. This observation leads to a conjecture that the VHE EGB is produced by a type of already known VHE \( \gamma \)-ray point sources with fluxes below the sensitivity of LAT. Since the dominant source class in the VHE band is that of BL Lac objects, a more precise conjecture is that the VHE EGB is produced by the unresolved BL Lac objects.

5. BL LAC OBJECT CONTRIBUTION TO VHE EGB

In the unification schemes of AGNs, BL Lac objects are identified with the FR I radio galaxies with jets aligned with the line of sight (Urry & Padovani 1995). This implies that the cosmological evolution of the BL Lac objects should follow that of the FR I radio galaxies. Recent studies of the cosmological evolution of FR I galaxies show that they experience “positive” cosmological evolution, which is usually described in terms of luminosity or comoving source density evolution as the increase of either average source luminosity or the average comoving source density with the redshift \( z \), as \((1 + z)^k , \ k > 0 \). Several recent studies find somewhat different values of \( k \), depending on the analyzed radio galaxy samples and different assumptions about the evolution type (luminosity or density), with \( k \) ranging in \( 1 \lesssim k \lesssim 3 \) (Sadler et al. 2007; Hodge et al.

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9 For the catalogs of extragalactic VHE \( \gamma \)-ray sources see, e.g., [http://tevcat.uchicago.edu](http://tevcat.uchicago.edu) and [http://www.isdc.unige.ch/vhe/index.html](http://www.isdc.unige.ch/vhe/index.html).
Since BL Lac objects are just the FR I galaxies specially oriented with respect to the line of sight, their cosmological evolution follows the evolution of FR I galaxies, with the increasing source luminosity or spatial density with the redshift. This implies that significant flux should be produced by the sources at large redshifts, \( z \sim 1 \). It is possible that most individual sources at large redshifts are too weak to be significantly detected by the LAT, but collective emission from all the set of BL Lac objects at high redshifts gives a significant contribution to the EGB.

Dependence of the total flux of the BL Lac population on the redshift could be found from the following straightforward calculation. Let us consider the total flux produced by sources at redshift \( z \) in a redshift interval \( \Delta z \). This redshift interval corresponds to the comoving distance interval \( \Delta r = \Delta z / H(z) \), where \( H(z) \sim \sqrt{\Omega_m + \Omega_b(1+z)^3} \) is the expansion rate of the universe filled with matter and cosmological constant with today’s densities \( \Omega_m \) and \( \Omega_b \).

As an example, we take the case of “pure luminosity” evolution with the average source luminosity increasing as \((1+z)^k\) and conserved comoving source density \( n(z) = n_0 = \text{constant} \). The number of sources in a spherical layer of thickness \( \Delta z \) is \( N(z) = 4\pi r^2 \Delta r \). Each source produces the flux in a given energy band \( F \sim (1+z)^{\Gamma k} / (4\pi r^2) \), where \( \Gamma \) is the photon index, and the factor \((1+z)^{\Gamma k}\) describes the change in the number of photons in a given energy band due to the cosmological redshift of the photon energies. One power of \((1+z)\) is compensated by the time delay between subsequent photons.

The flux from the sources at large redshifts is affected by the absorption of VHE photons on EBL. For example, at \( z \approx 1.5 \) the absorption modifies the source spectrum above the energy \( E \approx 50 \text{ GeV}, \) if one assumes the EBL evolution calculated by Franceschini et al. (2008). Absorption on EBL leads to suppression of the flux by a factor \( \exp(-\tau(E,z)) \), where \( \tau(E,z) \) is the optical depth with respect to the pair production.

The overall flux from the sources in the redshift interval \( \Delta z \) is

\[
\frac{\Delta F(E,z)}{\Delta z} = F \Delta N_s \sim \frac{(1+z)^{\Gamma k-1} e^{-\tau(E,z)}}{\sqrt{\Omega_m + \Omega_b(1+z)^3}}.
\]

Figure 5 shows the number of \( \gamma \)-rays as a function of source redshift. For this, we used BL Lac objects with known redshifts from the Veron and Veron (Veron-Cetty & Veron 2010) catalog complemented by BL Lac objects detected by the LAT, but not listed in the Veron and Veron catalog. Only sources with \( |b| > 10^\circ \) were considered. Here, we plot photon distributions from Fermi BL Lac objects in the three energy bands: 6.25–12.5 GeV, 25–50 GeV, and 100–200 GeV. One can see that at lower energies \( E < 50 \text{ GeV} \) a significant flux is produced by BL Lac objects at large redshifts up to \( z = 1.5 \). At the highest energies only the contribution from nearby sources at \( z < 0.7 \) is present. Two effects might explain the deficit of high-redshift sources at high energies. First, the flux at the highest energies is suppressed by absorption on EBL. Next, the photon statistics in the highest energy bin are low so that sources contributing to the flux in the 6.25–12.5 GeV bin produce less than one photon in the 100–200 GeV bin.

In the same figure, we also show the expected dependence of the number of photons on the redshift expected in different evolution models (Equation (4)). The models for cases \( k = 1, 2, 3 \) are shown with magenta lines for 6.25–12.5 GeV energy band. We normalize the models to the number of photons in the first redshift bin, in which we have the most complete knowledge of the BL Lac population.

From the comparison of the evolution models with the data, one might get an impression that the \((1+z)^3\) model is more consistent with the data than the models assuming faster evolution. However, the histogram in Figure 5 does not take into account photons from BL Lac objects with unknown redshifts. These BL Lac objects produce about 30% of all cumulative BL Lac flux. This means that at least 30% of the contribution to the overall flux (integrated over all redshifts) is missing in Figure 5. The model with evolution \( k = 1 \) predicts the total number of photons which is \( \sim 3\sigma \) below the total number of photons in BL Lac objects with known and unknown redshifts together in the 6.25–12.5 GeV energy band. In the energy band 3.125–6.25 GeV the underprediction of the total number of photons from BL Lac objects in the \( k = 1 \) model is at \( >5\sigma \) level, which means that the model is efficiently ruled out. Thus, Fermi-LAT observations of BL Lac objects indicate that BL Lac objects have positive cosmological evolution with \( k > 1 \).

In all other cases, \( k > 1 \), the discrepancy between the observed and expected number of photons from BL Lac objects starts already at small redshifts, \( z \gg 0.2 \). The “missing BL Lac” \( \gamma \)-rays could come either form BL Lac objects with unknown redshifts or from Fermi sources which are not yet identified as BL Lac objects or, finally, from BL Lac objects with fluxes below the sensitivity of the LAT.

Although individual high-redshift BL Lac objects would not be detectable by LAT, the cumulative flux of these BL Lac objects could give significant contribution to the EGB. Figure 6 shows the contributions from “missing BL Lac objects” at high redshifts up to \( z = 1 \) expected in four different models of cosmological evolution of BL Lac/FIR population. To calculate this contribution, we have normalized \( \Delta F/\Delta z \) distribution on the measured flux of BL Lac objects in the redshift bin \( 0 < z < 0.1 \). For the two last bins, 100–200 GeV and 200–400 GeV, the statistics of the BL Lac signal is too low to normalize \( \Delta F/\Delta z \) by the flux in the first redshift bin. To estimate the normalization of \( \Delta F/\Delta z \) in these energy bins, we have assumed that the average BL Lac spectrum extends as a power law with photon index \(-2\) up to the 400 GeV energy range. The resulting statistics of the signal in the last 200–400 bin in Figure 6 is low and is subject to large fluctuations.

From Figure 6 one could see that the conjecture that EGB is produced by collective emission from distant BL Lac objects is valid if BL Lac objects experience positive luminosity or density evolution of the form \((1+z)^k\) with \( k \approx 3 \). Slower or faster cosmological evolution with \( k = 2 \) or \( k = 4 \), respectively, would underproduce or overproduce the EGB. Since evolution \( k = 4 \) predicts number of photons more then diffuse gamma-ray background (see Figure 6) it is also excluded. Thus, the LAT data impose a constraint on the cosmological evolution of the BL Lac/FIR I population:

\[
1 < k < 4.
\]

From Figure 6 one could also see that if \( k = 3 \), all the flux of EGB above 30 GeV could be explained by a cumulative emission from the BL Lac population. At the same time, if \( k \) is significantly smaller than 3, as is suggested by some recent studies of the evolution of the parent population of FR I galaxies (Smolcic et al. 2009), significant part of the VHE EGB flux should come from a yet unknown source population or have a truly diffuse nature.

\( \gamma \)-ray flux from high-redshift BL Lac objects is modified in the VHE band by the effect of absorption on EBL. The
model spectra shown in Figure 6 take this effect into account. We use the model of Franceschini et al. (2008) to estimate the attenuation of the VHE $\gamma$-ray flux from redshifts up to $z = 1.5$. Model curves shown in Figure 6 assume that the average intrinsic spectrum of BL Lac objects does not have an HE cutoff up to $\sim 400$ GeV. The observed suppression of the flux above 100 GeV is explained only by the effect of absorption on EBL.

6. DISCUSSION AND CONCLUSIONS

In this paper, we have derived a measurement of EGB in the 10–400 GeV energy range from the analysis of Fermi/LAT

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Figure 5. Number of detected VHE photons as a function of redshift in the 6.25–12.5 GeV (blue dotted histogram), 25–50 GeV (green dashed histogram), and 100–200 GeV (red solid histogram) bands. Also shown are the distributions of photons with the redshift expected for different laws of BL Lac cosmological evolution. (A color version of this figure is available in the online journal.)

Figure 6. VHE EGB produced by unresolved BL Lac objects under different assumptions about the cosmological evolution of BL Lac population (the evolution law is marked to the left of each curve). The red data points show the EGB spectrum from Figure 3. (A color version of this figure is available in the online journal.)
data in the North and South Galactic Poles regions, $|b| > 60^\circ$. Our approach was to count all the protons detected by LAT in this region and estimate the number of counts from the point sources, from the Galactic diffuse emission, from the residual cosmic-ray background, and from the EGB. Subtracting the source, Galactic diffuse, and residual cosmic-ray background counts from the total number of counts in the North and South Galactic Pole regions we derived the EGB spectrum shown in Figure 3.

Comparing the spectrum of EGB in the $> 30$ GeV energy band with the spectrum of extragalactic point sources in the same energy band (Figure 4), we have noticed that the two spectra closely follow each other. Based on this observation, we have put forward a conjecture that the EGB above 30 GeV is explained by the unresolved BL Lac objects, which give a dominant contribution to the extragalactic point-source flux in this energy band. We have demonstrated that this conjecture is consistent with the EGB measurement provided that BL Lac objects follow positive cosmological evolution with the overall power of emission from the source population increasing as $(1 + z)^3$ up to $z \sim 1$ (Figure 6). Such cosmological evolution is roughly consistent with the measurements of cosmological evolution of FR I radio galaxies which are believed to be the parent population of BL Lac type objects and are also observed to have positive cosmological evolution of the form $(1 + z)^k$ with an uncertain value of $k$ between 1 and 3 (Rigby et al. 2008; Smolcic et al. 2009; Sadler et al. 2007). At the same time, it is opposite to the negative cosmological evolution of the HE-peaked BL Lac objects (Gioianni et al. 1999, 2001), which constitute a sub-class of the GeV–TeV $\gamma$-ray emitting BL Lac objects considered in our analysis.

If the VHE EGB is indeed produced by distant BL Lac objects at redshifts up to $z \sim 1$, the LAT will not be able to resolve it into point sources. Indeed, the brightest BL Lac object on the sky, Mrk 421, produced $\sim 30$ photons above 100 GeV. If the positive cosmological evolution of BL Lac population is mostly due to the increase of the comoving source density rather than an increase of the typical source luminosity, the brightest BL Lac objects at redshift $z \sim 1$ produce $\sim 10^{-5}$ $\gamma$-rays in LAT over some 2.5 years of exposure. This means that the LAT would not collect sufficient photon statistics to detect distant BL Lac objects individually. If the $(1 + z)^3$ evolution is mostly due to the increase of the average source luminosity, BL Lac objects at redshift 1 have a luminosity an order-of-magnitude higher than local BL Lac objects. However, even with higher luminosity, they produce on average 0.1 photon in LAT, so that they are still not individually detectable.

If the real value of $k$ is much below $k = 3$, as indicated in a recent study by Smolcic et al. (2009), emission from the unresolved BL Lac objects would not explain the VHE EGB flux and there should be another source class or a mechanism of production of diffuse emission which would account for the EGB. Such a mechanism might, in fact, be indirectly related to the BL Lac population. Most of the power output from BL Lac objects at the energies above 100 GeV is converted into the electromagnetic emission from $\gamma$-ray-induced cascade in intergalactic medium. The intrinsic spectra of BL Lac objects (and of FR I radio galaxies, such as M87 and Cen A; Aharonian et al. 2006, 2009) are known extend up to $\sim 10$ TeV. Typical energy of the cascade photons which are produced via IC scattering of CMB photons by the $e^+ e^-$ pairs deposited in the intergalactic medium is $E_\gamma \sim 100 E_\gamma / 10$ TeV$^2$ GeV (Neronov & Semikoz 2009). If the intrinsic source luminosities in the $10^\circ \rightarrow 10$ TeV range are comparable to the luminosities in the 10–100 GeV, the total flux of the cascade emission in the 10–100 GeV band is expected to be comparable to the point-source flux, so that the cascade emission could give a significant contribution to the EGB (Coppi & Aharonian 1997). This is consistent with the observation that the VHE EGB level is comparable to the cumulative extragalactic point-source flux in the 10–100 GeV band observed by LAT.

The only possibility to test the hypothesis of BL Lac origin of EGB would be to use deep observations with ground-based $\gamma$-ray telescopes. Ground-based $\gamma$-ray telescopes, which are sensitive in the VHE energy band, have much larger collection area above several hundreds of GeV and, as a consequence, could detect much weaker sources than LAT. The flux from Mrk 421-like BL Lac objects at the redshift $z \sim 1$ is $\sim 10^{-13}$ erg cm$^{-2}$ s$^{-1}$. If the cosmological evolution of BL Lac objects is due to an increase of the average source luminosity with redshift, the brightest BL Lac objects at redshift $z \sim 1$ might produce fluxes up to $10^{-12}$ erg cm$^{-2}$ s$^{-1}$ at 100 GeV. At the energies around $E \lesssim 100$ GeV this flux is not strongly attenuated by the absorption on EBL. The energy $E \gtrsim 100$ GeV is around or below the low energy threshold of the current generation Cherenkov telescopes, such as HESS, MAGIC, and VERITAS. However, next generation facilities, Cherenkov Telescope Array (CTA; CTA Consortium 2010) or 5@5 (Aharonian et al. 2001) are expected to have an energy threshold significantly below 100 GeV. Their sensitivity could be sufficient to resolve the VHE EGB into point sources, at least in the case when the cosmological evolution is mostly luminosity, rather than source density evolution.

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REFERENCES

Abdo, A. A., Ackermann, M., Ajello, M., et al. 2010a, Phys. Rev. Lett., 104, 101101
Abdo, A. A., Ackermann, M., Ajello, M., et al. 2010b, ApJ, 720, 435
Abdo, A. A., Ackermann, M., Ajello, M., et al. 2010c, J. Cosmol. Astropart. Phys., JCAP04(2010)014
Abdo, A. A., Ackermann, M., Ajello, M., et al. 2010d, ApJS, 188, 405
Ackermann, M., Ajello, M., Atwood, W. B., et al. 2010, Phys. Rev. D, 82, 092004
Aharonian, F. A. (ed.) 2004, Very High Energy Cosmic Gamma Radiation (River Edge, NJ: World Scientific)
Aharonian, F. A., Akhperjanian, A. G., Anton, G., et al. 2009, ApJ, 695, L40
Aharonian, F. A., Akhperjanian, A. G., Barres de Almeida, U., et al. 2008, Phys. Rev. Lett., 101, 261104
Aharonian, F. A., Akhperjanian, A. G., Bazer-Bachi, A. R., et al. 2006, Science, 314, 1424
Aharonian, F. A., Konopelko, A. K., Völk, H. J., & Quintana, H. 2001, Astropart. Phys., 15, 335
Atwood, W. B., Abdo, A. A., Ackermann, M., et al. 2009, ApJ, 697, 1071
Bereczky, V., Gazazov, A., Kachelriess, M., & Ostapchenko, S. 2011, Phys. Lett. B, 695, 13
Bereczky, V. S., & Smirnov, A. Yu. 1975, Ap&SS, 32, 461
Chiang, J., Fichtel, C. E., von Montigny, C., Nolan, P. L., & Petrosian, V. 1995, ApJ, 452, 156
Coppi, P., & Aharonian, F. 1997, ApJ, 487, L9
CTA Consortium 2010, arXiv:1008.3703
Fichtel, C. E., Simpson, G. A., & Thompson, D. J. 1978, ApJ, 222, 833
Franceschini, A., Rodighiero, G., & Vaccari, M. 2008, A&A, 487, 837
Gilmore, R. C., Madau, P., Primack, J. R., Somerville, R. S., & Haardt, F. 2009, MNRAS, 399, 1694
Gioianni, P., Menna, M. P., & Padovani, P. 1999, MNRAS, 310, 465
