NEAR-INFRARED ADAPTIVE OPTICS IMAGING OF INFRARED LUMINOUS GALAXIES: THE BRIGHTEST CLUSTER MAGNITUDE–STAR FORMATION RATE RELATION

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ABSTRACT

We have established a relation between the brightest super star cluster (SSC) magnitude in a galaxy and the host star formation rate (SFR) for the first time in the near-infrared (NIR). The data come from a statistical sample of \textasciitilde 40 luminous IR galaxies (LIRGs) and starbursts utilizing $K$-band adaptive optics imaging. While expanding the observed relation to longer wavelengths, less affected by extinction effects, it also pushes to higher SFRs. The relation we find, $M_K \sim -2.6 \log SFR$, is similar to that derived previously in the optical and at lower SFRs. It does not, however, fit the optical relation with a single optical to NIR color conversion, suggesting systematic extinction and/or age effects. While the relation is broadly consistent with a size-of-sample explanation, we argue physical reasons for the relation are likely as well. In particular, the scatter in the relation is smaller than expected from pure random sampling strongly suggesting physical constraints. We also derive a quantifiable relation tying together cluster-internal effects and host SFR properties to possibly explain the observed brightest SSC magnitude versus SFR dependency.

Key words: galaxies: star clusters: general – galaxies: star formation – infrared: galaxies

1. INTRODUCTION

Young massive star clusters also known as super star clusters (SSCs) are a good tracer of the current star formation of the host galaxy (Portegies Zwart et al. 2010). Over the last decade, many studies have shown that there is an empirical relation between the $V$-band luminosity of the brightest cluster magnitude\textsuperscript{7} and the global star-formation rate (SFR) of the galaxy (e.g., Larsen 2002; Weidner et al. 2004; Bastian 2008). Although various reasons have been suggested to explain the relation, they all highlight the important role star clusters play in understanding their host galaxy properties. Larsen (2002) and Whitmore (2003) emphasize the importance of size-of-sample effect in universal cluster formation, i.e., that large SSC populations preferentially sample the initial luminosity functions (LFs) to higher values, while Weidner et al. (2004) and Bastian (2008) show through theoretical simulations that effects from physical processes could also generate the observed correlation. Weidner et al. (2004) derive an expression which directly relates the total SFR of the galaxy with the mass of the brightest star cluster, with the assumption that the most massive cluster is always the brightest.

On the other hand, Bastian (2008) argue that the youngest clusters are the brightest implying the tight observed relation is an imprint of the current SFR of the galaxy. Adamo et al. (2011) find that clusters in blue compact galaxies (BCGs) preferentially lie above the relation fit to more “normal” galaxies suggesting that the environments of SSCs play a role in determining the relation.

So far the validity of the relation has only been tested in the optical regime, and mostly using fairly nearby star-forming galaxies. Extinction effects are difficult to correct for especially in the duster galaxies and will necessarily introduce scatter in the relation; using redder wavelengths will significantly improve

\textsuperscript{7} Hereafter, we will refer to the “brightest cluster magnitude” as the “brightest cluster.”

2. OBSERVATIONS AND PHOTOMETRY

Most of our sample, \textasciitilde 30 galaxies, were imaged in the $K_S$-band with the NACO AO-instrument on the ESO Very Large Telescope (VLT).\textsuperscript{8} The targets are IRAS galaxies from the RBGS (Sanders et al. 2003), selected to be $D_L \lesssim 200$ Mpc with IR luminosities above log $(L_{IR}/L_\odot) = 10.6$, and with a bright enough reference star near the field-of-view required by the AO system. The galaxies are hence a representative statistical sample of all IR-bright galaxies above log $(L_{IR}/L_\odot) = 10.6$ within the distance limit. They all have “cool” IRAS colors, though no active galactic nuclei were excluded a priori. Depending on the size of the galaxy, either the S27 or S54 camera was used, resulting in a pixel scale of 0.027 pixel$^{-1}$ or 0.054 pixel$^{-1}$, respectively. The final point spread function (PSF) resolution was typically \textasciitilde 0.1. We used a dithering mode with 120 s per pointing, with total integration times per target ranging between 20 and 40 minutes.

Our IRAF-based pipeline was used to perform sky-subtraction. The individual frames were then aligned before average-combining them to get the final science image. We checked for image quality in individual frames and those with non-optimal AO-corrections were excluded, resulting in shorter total integration times in some cases (Table 1). Some frames with obviously non-photometric conditions were also excluded.

\textsuperscript{8} Programs 086.B-0901 (PI: Escala) and 089.D-0847 (PI: Mattila).
Photometric zero-points were either retrieved from the ESO/NACO official Web site, if recorded, or estimated by correlating with Two Micron All Sky Survey $K_S$ point-sources present in the field. Fields where both methods were available were used to check consistency of photometry. The zero-point uncertainty of $\sim 0.1$ is included in resulting SSC photometry, which varies in the range $\sim 0.1$–$0.4$ mag. The results are in a Vega-based system.

We added 12 other luminous IR galaxies (LIRGs) which form a homogeneous data set (same wavelength and similar depth and resolution) with the new sample described above. Eight of these come from Gemini-N/ALTAIR/NIRI and four from earlier NACO data. The Gemini sample was selected with a higher log($L_{IR}/L_\odot$) $>$11.3 cut-off; the final results are checked without these in case of any biases. Ten of these additional targets were published in Randriamanakoto et al. (2013) with the

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Table 1
Observation log of the VLT/NACO and Gemini/NIRI $K$-band Data

| Galaxy Name | Exp. Time | $D_L$ | $M_{K}^{\text{exposure}}$ | log $L_{IR}$ | SFR |
|-------------|-----------|-------|-----------------|-------------|-----|
|             | (s)       | (Mpc) | (mag)           | ($L_\odot$) | (M$_\odot$ yr$^{-1}$) |
| VLT/NACO data |
| ESO 440–IG058-N | 2280 | 102.0 | −15.89 | 10.59 | 6.6 |
| IC 2522 | 1260 | 46.1 | −15.52 | 10.63 | 7.3 |
| MCG –02-01-052 | 1860 | 110.0 | −17.21 | 10.63 | 7.3 |
| NGC 3620 | 1780 | 24.9 | −15.65 | 10.70 | 8.5 |
| ESO 428–G023 | 1770 | 44.5 | −15.51 | 10.76 | 9.8 |
| ESO 221–IG008 | 4620 | 46.7 | −15.90 | 10.77 | 10.0 |
| IRAS 06164 + 0311 | 1100 | 41.5 | −16.06 | 10.79 | 10.5 |
| NGC 1134 | 1240 | 47.4 | −15.86 | 10.83 | 11.5 |
| ESO 491–G020 | 1200 | 43.5 | −17.47 | 10.86 | 12.3 |
| NGC 4433 | 1020 | 46.3 | −16.39 | 10.87 | 12.6 |
| NGC 1204 | 1100 | 61.4 | −17.50 | 10.88 | 12.9 |
| NGC 3508 | 1860 | 59.1 | −16.42 | 10.90 | 13.5 |
| NGC 1819 | 1950 | 61.9 | −17.11 | 10.90 | 13.5 |
| MCG –02-33-098 | 960 | 70.8 | −15.69 | 10.95 | 15.3 |
| NGC 4575 | 720 | 45.0 | −15.62 | 10.96 | 15.5 |
| NGC 6000 | 2880 | 32.1 | −16.36 | 10.97 | 15.8 |
| ESO 550–IG025-S | 1950 | 135.0 | −16.76 | 11.03 | 18.2 |
| ESO 319–G020 | 960 | 43.2 | −15.98 | 11.04 | 18.6 |
| MCG +02-20-003 | 2340 | 70.5 | −16.19 | 11.08 | 20.4 |
| ESO 264–G057 | 1120 | 75.8 | −17.15 | 11.08 | 20.4 |
| ESO 320–G030 | 1120 | 49.0 | −15.70 | 11.10 | 21.4 |
| ESO 221–IG010 | 240 | 45.9 | −17.19 | 11.17 | 25.1 |
| ESO 267–G030 | 1920 | 80.9 | −16.86 | 11.19 | 26.3 |
| ESO 550–IG025-N | 1950 | 135.0 | −17.07 | 11.24 | 29.5 |
| ESO 440–IG058-S | 2280 | 102.0 | −17.79 | 11.28 | 32.3 |
| NGC 3110 | 2520 | 75.2 | −17.74 | 11.31 | 34.7 |
| IRAS 13052–5711 | 3480 | 91.6 | −16.51 | 11.34 | 37.2 |
| ESO 264–G036 | 2340 | 92.0 | −17.98 | 11.35 | 38.1 |
| IRAS 12116–5615 | 1800 | 117.0 | −18.42 | 11.59 | 66.1 |
| IRAS F06076–2139 | 1560 | 160.0 | −19.61 | 11.59 | 66.1 |
| IRAS 01173+1405 | 1320 | 127.0 | −18.00 | 11.63 | 72.5 |
| IRAS 01364–1042 | 1560 | 201.0 | −20.01 | 11.76 | 97.8 |
| IRAS 18293–3413 | 1230 | 74.6 | −18.23 | 11.81 | 109.8 |
| NGC 6240 | 720 | 103.0 | −18.79 | 11.85 | 120.4 |
| IRAS 19115–2124 | 1410 | 206.0 | −19.71 | 11.87 | 126.0 |

| Gemini/NIRI data |
| IRAS F16516–0948 | 900 | 94.8 | −18.10 | 11.24 | 29.5 |
| CGCG 049–057 | 1680 | 56.4 | −17.14 | 11.27 | 31.6 |
| IRAS F17578–0400 | 1470 | 57.3 | −17.56 | 11.35 | 38.1 |
| MCG +08-11-002 | 1140 | 79.9 | −17.52 | 11.41 | 43.7 |
| IRAS F17138–1017 | 990 | 72.2 | −18.48 | 11.42 | 44.7 |
| NGC 3690 | 2192 | 45.3 | −17.88 | 11.48 | 51.6 |
| IC 694 | 1260 | 45.3 | −16.95 | 11.66 | 77.4 |
| IC 883 | 1440 | 101.0 | −18.38 | 11.67 | 79.5 |

Notes. Column 1: IRAS survey name; Column 2: total exposure time; Column 3: luminosity distance from NED Database; Column 4: $K$-band absolute magnitude of the brightest cluster; Column 5: galaxy IR luminosity from Sanders et al. (2003), any value marked by † is estimated by using the method described in Section 3; Column 6: SFR derived from Equation (1).
SSC catalogs ready for analysis, and two unpublished targets are analyzed here (NGC 6240 and NGC 6000). We refer the reader to Mattila et al. (2007), Kankare et al. (2008, 2012), Väisänen et al. (2008a, 2008b), and Randriamanakoto et al. (2013) for details of observations and data reduction.

2.1. Photometry and SSC Selection

The SSC candidate photometric catalogs were generated by following the same steps as in Randriamanakoto et al. (2013). Briefly, an unsharp-masked version of the science image was used for object detection using SExtractor (Bertin & Arnouts 1996) with a configuration optimized to detect faint sources in very complex backgrounds. Photometry itself was then performed with IRAF/PHOT in aperture radii of 2 and 3 pixels (0′′11 and 0′′08) in frames taken with SS4 and S27, respectively. The sky annuli were 1.5 and 2 pixels wide, respectively, with the inner radius one pixel away from the aperture radius in both cases. Aperture correction was achieved with the usual curve-of-growth method, drawn until ~1′′ radius. If there were enough isolated point sources in the field, this aperture correction was dependent on the distance to the AO reference star. Our SSC selection method retains only likely point-sources by utilizing the concentration of light in the detected sources. Simultaneously, we exclude likely foreground stars using visual inspection supported by expected star counts from the models of Besançon (Robinet et al. 2003), as well as the nucleus, or nuclei, of the target galaxies.

3. RESULTS

Assuming, as usual, that the SFR of strongly star-forming galaxies is well represented by their IR luminosity, the following empirical relation by Kennicutt (1998) was used to convert the \( L_{\text{IR}} \) to SFR for each galaxy:

\[
\frac{\text{SFR}}{M_\odot \text{yr}^{-1}} = 1.7 \times 10^{-10} L_{\text{IR}} (L_\odot).
\]  

(1)

Both values are listed in Table 1.

If the galaxy is a close pair or a multiple system, its IR luminosity has to be separated into individual components since the Sanders et al. (2003) values are for the whole system due to the poor spatial resolution of IRAS. Therefore, we used WISE 12 and 22 μm, Spitzer/MIPS 24 and 60 μm, and Herschel/PACS 70 μm archival data\(^{10}\) to measure the flux from each galaxy component and the IRAS-based \( L_{\text{IR}} \) was redistributed according to the average ratio from all those wavelengths that were available and resolved. This method was adopted for these targets: ESO 440−IG058, MCG−02-01-052, MCG−02-33-098, ESO 491−IG020, ESO 550−IG025, IC 694/NGC 3690 (=Arp 299).

The NIR brightest cluster–SFR relation is shown in Figure 1. A weighted linear fit to all the points (the uncertainties of the resolved components’ \( L_{\text{IR}} \) points were doubled) results in the following relation, shown as the dashed line in Figure 1:

\[
M_K^{\text{brightest}} = -3.10 \times \log \text{SFR} - 12.75.
\]  

(2)

Blending effects may be a concern, however. The physical spatial resolution in our survey, corresponding to the ~0′′1 PSF size, is typically 20–60 pc. Thus, individual detections of SSC candidates could potentially be blends of more than one intrinsic SSC. The effect may not be too severe for the brightest clusters, since Randriamanakoto et al. (2013) showed using Monte Carlo (MC) simulations that a single bright SSC will overwhelmingly dominate the luminosity of a SSC candidate detection when small apertures are used, except in the most distant targets approaching \( D_L \sim 200 \text{ Mpc} \), or in case of very strong clustering of SSC regions. Nevertheless, we check for blending effects in the following way: the \( M_K \) versus SFR relation was fit for the “safe” galaxies at \( D_L \leq 80 \text{ Mpc} \) (the slope is ~2.50 in this case), and Figure 2 then plots the difference of the brightest \( M_K \) from this best-fit relation versus the distance of the host. If distance plays no part, a scatter plot is expected. Indeed, no systematics are seen, apart from the three most distant targets falling significantly above the null-hypotheses line. We interpret this as the brightest SSC in those three targets potentially being

\(^{10}\) http://irsa.ipac.caltech.edu, http://herschel.esac.esa.int
contaminated by other clusters and exclude them from further analysis. We obtain a new best fit (solid line in Figure 1) when using a $D_L \leq 150$ Mpc constraint:

$$M_K^{\text{brightest}} = -2.56 \times \log \text{SFR} - 13.39.$$  \hfill (3)

The slope varies in the range $-2.49$ to $-2.56$ with distance limits set in between 80 and 150 Mpc. The formal uncertainty of the slope is $\pm 0.07$ and the scatter of the observed $M_K$ values is $\sigma \approx 0.62$ mag. The slope without the eight Gemini galaxies is $-2.43$; the difference is within $\approx 1.5\sigma$ of the overall fit and does not affect any conclusions.

4. DISCUSSION

4.1. Statistical Interpretation

It has been suggested that the $M^{\text{brightest}}$–SFR relation in the V-band can be explained in purely statistical terms (Larsen 2002). To explore this possibility in the NIR, $M_K^{\text{brightest}}$ is plotted against the number of SSC candidates brighter than a certain absolute magnitude level. We select $M_K = -15$ mag since at that level we do not yet need completeness corrections. The result can be seen in Figure 3. This correlation is consistent with the idea that the more clusters are forming in a galaxy, the higher will be the probability to sample the brightest ones from a given LF of the overall population (Larsen 2002). The empirical relation can also be tied to the slope of the LF. If the $L_{\text{max}}$ of the most luminous object scales with the total number $N$ of the clusters as given by Whitmore (2003)

$$L_{\text{max}} \sim N^{\eta},$$  \hfill (4)

then by using the equation from Hunter et al. (2003)

$$\eta = \frac{1}{\alpha - 1},$$  \hfill (5)

we can derive the power-law slope $\alpha$ of the cluster LF at higher luminosities. The best-fit $\eta = 0.64$ corresponds to $\alpha = 2.56$ which is not unreasonable for bright parts of SSC LFs (Portegies Zwart et al. 2010). From the 10 LIRGs in Randriamanakoto et al. (2010) we found a slope of $\alpha \sim 1.9$ (flatter than for normal spirals) which corresponds to $\eta = 1.1$—this is overplotted in Figure 3 as the dashed line and is seen to represent the data at $\log N > 0.5$ fairly well.

The scatter of $\sigma \approx 0.62$ mag in Figure 1 is smaller than in relations derived in the optical showing $\sigma \sim 1$ mag (Larsen 2002), unsurprising since extinction effects are smaller. Given that the SFR determination uncertainties should be of the order of 0.4 dex, the observed scatter is surprisingly small, however. We ran MC simulations where a given LF with no physical upper limit was sampled purely randomly and the magnitude of the brightest cluster was recorded. The $M^{\text{brightest}}$ distribution is narrower with steeper LFs parameterized with the power-law index $\alpha$. In particular, we find $\sigma = 1.15$ for $\alpha = 2.0$, and $\sigma = 0.77$ and $0.58$ for $\alpha = 2.5$ and 3.0, respectively, each with uncertainties of $\approx 0.02$. In case of purely statistical sampling the scatter in our relation must hence be the result of steeper than observed LFs of $\alpha > 2.5$—or there are other physical characteristics at play which determine the luminosity of the brightest cluster.

In summary, while the characteristics of the number of detected SSCs versus SFR are consistent with a size-of-sample effect, the tightness of the brightest cluster versus SFR relation in particular suggests that it would be premature to reject an underlying physical cause for this relation.

![Figure 3. $M_K^{\text{brightest}}$ as a function of the number of SSCs brighter than $-15$ mag for targets with $D_L \leq 150$ Mpc. The solid line is the best-fit relation, whereas the dashed line corresponds to a LF power-law slope of $\alpha = 1.9$. Symbols as in Figure 1.](image)

4.2. Physical Interpretation

Clusters are born of collapsing giant molecular clouds which inevitably are affected by their galactic environments, especially in cases of interacting and merging galaxies—how this environment exactly defines SSC properties, and disruption, is a matter of intense debate (e.g., Lamers 2009). Specifically regarding the brightest cluster versus SFR relation, Adamo et al. (2011) have shown how SSCs in BCGs appear elevated from the general relation, and suggest this could be a result of a higher cluster formation efficiency in their extreme environments. Larsen (2009) and Gieles (2009) suggest that the characteristic cluster mass may change as a function of environment, and grow in the more intense SFR of interactions and mergers. One appealing possibility for such a change is the lack of large scale rotation in galaxy mergers (Escala & Larson 2008; Weidner et al. 2010; Escala et al. 2013).

In addition, there might well be internal constraints on SSC properties. In the following we outline a possible physical interpretation of the brightness versus SFR relation based on the idea that the total luminosity, and mass, of a stellar cluster is weighted toward its highest mass stars, and these stellar masses may also be correlated with the environments of the clusters (e.g., Weidner et al. 2009).

The total luminosity of a cluster can be computed for a given initial mass function (IMF) and mass–luminosity relation. Assuming a mass–luminosity relation of the form $L \propto m^\gamma$ and a power-law IMF ($dN/dm \propto m^{-\beta}$), the total luminosity of a cluster is given by

$$L_{\text{tot}} \propto M_{\text{cl}}^{(\alpha-\beta+1)\gamma},$$  \hfill (6)

where $M_{\text{cl}}$ is the total mass of the cluster, which is assumed to satisfy a relation with the most massive star of such cluster: $M_{\text{star}}^{\text{max}} \approx M_{\text{cl}}^{\gamma}$, being $\gamma \approx 0.45$ estimated from observations and $\gamma \approx 2/3$ predicted from simulations (see Weidner et al. 2009 for a review on different estimates).

Assuming that the most massive unstable gas cloud in a galaxy, $M_{\text{cloud}}^{\text{max}} \approx f_g^{1/2} M_{\text{gas}}$ (Escala & Larson 2008), leads to
the formation of the most massive SSC ($M_\mathrm{cl} \propto M_\mathrm{max}^{\alpha_1}$) and taking into account the correlation between such a cloud and the SFR in galaxies, $SFR \propto [M_\mathrm{max}^{\beta}]^\delta$ with $\delta \sim 1.5$ (Escala 2009, 2011), the total luminosity of the brightest cluster is given by $L_{\mathrm{tot}}^{\mathrm{brightest}} \propto SFR^{(\alpha_1 - \beta + 1)\gamma}/\delta$. Finally, this can be expressed in terms of absolute magnitude by $M_K^{\mathrm{brightest}} \propto -2.5 \log L_{K}^{\mathrm{brightest}}$, resulting in:

$$M_K^{\mathrm{brightest}} \propto -2.5(\alpha_1 - \beta + 1)\gamma/\delta \log \text{SFR}.$$  

For a Salpeter IMF ($\beta = 2.35$), $\delta = 1.5$ (Escala 2009, 2011), $\gamma \sim 0.45$ estimated from observations (Weidner et al. 2009) and a slope of the mass–luminosity relation of $\alpha_1 \sim 5$, Equation (7) gives a slope closer to $-3$ in the brightest cluster–SFR relation, which is comparable to the slope observed if we use the whole sample (Equation (2)). On the other hand, for a slope of the mass–luminosity relation of $\alpha_1 \sim 4$, Equation (7) gives a slope $\sim -2$, which is closer to that derived from our data excluding potentially blended cases (Equation (3)).

Unfortunately, we do not have a good estimate for the slope of the mass–luminosity relation $\alpha_1$ at high masses—it may, for example, vary in between 1.76 and 8.87 depending on the highest mass of a star in a cluster for masses larger than 7 $M_{\odot}$ (Parravano et al. 2003). Nevertheless, Equation (7) may be used as a simple physical interpretation of the effects characterizing SSC properties at scales ranging from internal to galactic.

### 4.3. Comparison to the $V$-band Relation

Figure 4 shows the brightest cluster–SFR relation with an expanded scale. The triangles are $V$-band data as compiled by Adamo et al. (2011), assuming a constant $V - K = 2$ conversion, typical for a ~10 Myr age stellar population. The solid line is the best-fit slope of $-2.56$ from our own NIR data extrapolated to lower SFRs. It appears that the optical points would require a slightly flatter slope, and indeed Weidner et al. (2004) find $\sim -1.9$ shown as the dashed line, though Larsen (2002) derives $\sim -2.5$ from a subset of the data. The simplest explanation could perhaps be extinction: the highest SFR galaxies, LIRGs and ultraluminous IR galaxies, are predominantly interactions and mergers with more dust on average than lower SFR galaxies (Piquéras López et al. 2013). The uncorrected optical points at higher SFR could lie too low artificially, thus flattening the slope. However, with this data set alone it is not possible to confirm this—it could as well be that the brightest clusters detected in $V$-band are not necessarily (always) the most luminous clusters in $K$-band possibly implying age effects. That the slope appears slightly different necessarily points to some systematic effects along the SFR base-line, i.e., a constant $V - K$ shift is not appropriate.

The slope may become steeper if the points at higher SFR, preferentially, are bound results of mergers of individual SSCs in very dense star cluster complexes (Fellhauer & Kroupa 2005). We also note that outliers below the relation can be understood as cases where the brightest SSC is not detected, or is severely extincted. As discussed by Bastian (2008), outliers above the line may be cases where the detected cluster is significantly older than the general population of brightest SSCs. These are very interesting questions to tackle with combinations of optical and NIR data in the future.

### 5. SUMMARY AND FUTURE WORK

From a $K$-band AO sample of 43 strongly star-forming galaxies, mostly LIRGs, we establish the brightest cluster magnitude–SFR relation in the NIR regime. The relation is much less affected by extinction effects than similar comparisons in the $V$-band. We find a slope of $-2.56$ which is similar to those from optical derivations made at lower SFR levels, though the extension of our slope appears to not be consistent with the full range of optical SSC luminosities if a single $V - K$ conversion is adopted. We suggest that a systematic extinction effect, where SSCs in higher SFR hosts live in dustier environments, would be a simple explanation for the trend, but systematic age differences may also be involved.

A good correlation of the most luminous cluster and the number of SSCs with $M_K$ magnitude brighter than $-15$ shows that a size-of-sample effect is broadly consistent with the observed $M_K^{\mathrm{brightest}}$–SFR relation. On the other hand, the observed scatter in the relation is surprisingly small, and we show that it can be explained with random sampling effects only if the LF of SSCs is very steep at the bright end, steeper than usually observed. Hence, physical reasons determining the luminosity of the brightest SSC from host properties, and/or internal cluster effects, likely play a role as well. We derived a relation tying the stellar IMF and mass–luminosity relations together with the global SF properties of the host in explaining the observed brightest cluster magnitude–SFR relation.

In the next steps of the work, we will investigate in more detail the environments and extinctions of the host galaxies, and masses and ages of the SSCs inside them, with a combination of optical and NIR data and kinematic information. These will allow more secure disentanglement of the various physical effects governing the lives and characteristics of SSCs in galaxies.

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