Reproducing laboratory earthquakes with a discrete-continuum model

Guilhem Mollon, Jérôme Aubry, and Alexandre Schubnel

1 Université de Lyon, LaMCoS, INSA-Lyon, CNRS UMR5259, F-69621, France
2 Laboratoire de Géologie, École Normale Supérieure/CNRS UMR 8538, PSL Research University, F-75005 Paris, France

Abstract. We present a novel numerical model allowing to take the best part of continuum-based and discrete modelling in a single framework. This model is applied to the reproduction of laboratory earthquakes in a high-pressure triaxial cell. It allows to represent most of the relevant phenomena at stake, including elastic stress build-up during loading, fast and slow sliding events, seismic waves emission in the surrounding elastic medium and evolution of fault gouge on the sliding interface. We review here some illustrative results obtained with this model and propose further research avenues.

1 Introduction

Laboratory earthquakes have become a major tool for the understanding of seismic and aseismic sliding on natural faults [1, 2, and references therein]. Outstanding progress has been made in the instrumentation of such experiments, but for practical reasons, it remains difficult for experimentalists to acquire the full mechanical and kinematic processes at stake on faults, at high sampling rates in space and time. This is where numerical simulations can be helpful.

We present a numerical model which aims at reproducing triaxial experiments on saw-cut marble samples described in [1, 2]. This model is in 2D and miniaturized when compared to its real-world counterpart, but includes relevant deformation mechanisms. The hanging wall and footwall of the sample are represented using a meshfree continuum approach, apart from the areas located in the neighbourhood (within a few hundreds of micrometers) of their contacting surfaces. In these areas, conforming polygonal grains bonded by a cohesive-zone model are implemented in a DEM framework [3]. Realistic boundary conditions (in terms of the elasticity of the loading system, of the absorption of the elastic waves and of the fluid pressure applied on the lateral boundaries) are introduced. Constitutive laws (both in the continuum and discrete parts) are calibrated based on experimental results found in the literature [4].

Figure 1. Model setup, with continuum and discrete areas.
2 Simulations

As for the experiments reported in [1, 2], the numerical model reproduces the deformation of a saw-cut sample (Fig. 1). The sample is 8 mm-high and 4 mm-wide (1/10 scale with respect to the reference experiment). Each part of the fault is continuous and elastic, excepting the first 200 µm of the contacting surfaces which are composed of discrete rigid and conforming polygons (typical size of 10 µm). They interact through a Cohesive Zone Model (CZM), which means that they reproduce well the elastic properties of the bulk rock but can also be damaged and separated from the surface if a threshold shear or tensile strength is overcome at their boundaries. The CZM parameters are calibrated using additional simulations reproducing published results of triaxial compression tests on marble [4] (Fig. 2).

Figure 2. Independent Calibration of the Cohesive Zone Model parameters based on experiments reported in [4]. Low confining stress leads to localized and brittle failure while large confining stress promotes distributed damage

The saw-cut sample is submitted to a lateral confining stress $\sigma_3$ (set to 45, 90, or 180 MPa) and to vertical compression at constant speed, through two elastic blocks which are designed to mimic the deformability of the experimental apparatus and to dissipate elastic waves triggered by seismic events at the precut interface. Simulations are stopped for a vertical strain of 2%. The fault roughness is limited to the irregularities brought by the polygonal shapes of the rigid grains. The simulations are performed with the code MELODY [5], which allows mechanical coupling between large collections of rigid grains and continuous deformable bodies.

Figure 3. Shear stress field (projected along the fault); (a.) before deviatoric loading; (b.) before and (c.) after the first stress drop; (d.) at the end of simulation ($\sigma_3$=180 MPa)

3 Stress history

As shown in Fig. 3, simulations are first characterized by the progressive increase of stress fields within the samples, combined with damage on the first grain layers at the interface. At a critical differential stress, sliding events occur and stress values drop and become more heterogeneous. Time-series of stress and strain tensors show typical laboratory earthquake cycles, illuminating successive accumulation and release of shear stress on the fault (Fig. 4). The case $\sigma_3$=45 MPa leads to a large number of low-amplitude stress drops, while the case $\sigma_3$=180 MPa is characterized by a smaller number of large stress drops. The case $\sigma_3$=90 MPa shows an intermediate behaviour.

Figure 4. Shear stress history on the fault

When expressed in terms of friction coefficient on the fault (Fig. 5), numerical results reveal that both peak and plateau friction coefficients decrease with increasing confining pressure. For $\sigma_3$=45 MPa, friction coefficient reaches 0.40 before the first stress drop and stabilizes around 0.33 during established seismic cycle; While for $\sigma_3$=180 MPa, peak and average plateau friction coefficients are close to 0.32 and 0.25, respectively. This low friction is close to that observed in similar experimental conditions [1, 2].

Figure 5. Evolution of fault friction as a function of vertical shortening of the rock sample
4 Sliding velocity history

Fault sliding is recorded for each simulation at a high sampling rate, both in space and time. This is performed by monitoring displacements and velocities of the surface nodes of the continuum parts of the two half blocks and by projecting them on the mean fault plane. It provides a detailed velocity map of the fault for the three confining pressures (Fig. 6). A large number of events can be observed and characterized. Dedicated procedures are used for detection and characterization of each sliding event. To illustrate this database, we present here three characteristic events of increasing size. All of them occurred under a confining stress $\sigma_3=180$ MPa. The first illustrative event (Fig. 7a) is of small amplitude, with a stress drop of $\sim 0.4$ MPa and a sliding distance of $\sim 0.3$ µm. This event originates in the middle part of the fault and propagates towards its extremities.

Figure 6. Zoomed view on a small portion of the space-time maps of the sliding velocity on the fault, for the three confining stresses.

The second illustrative event (Fig. 7b) is of intermediate amplitude, with a stress drop of $\sim 1.3$ MPa and a sliding distance of $\sim 1$ µm. This event originates at the upper extremity of the fault and propagates towards the other extremity, where it triggers a secondary (and more intense) event which back-propagates towards the upper extremity. The third illustrative event (Fig. 7c) is of very large amplitude, with a stress drop of $\sim 16$ MPa and a sliding distance of $\sim 10$ µm. This fairly complex event includes a succession of smaller slip events which trigger each other over the whole fault. The main event occurs at a sliding velocity in the order of 10 m/s. We observe a major reorganization of the stress field in the sample during this event.

5 Fault interface damage evolution

Fig. 8 summarizes the evolution of fault surface damage in the discrete degradable part of the model. This damage increases progressively at each contact between neighbouring cells, and is averaged at each simulation step as a function to the distance to the main fault plane. The upper part of Fig. 8 shows that the damage spreads out during each simulation, both as the result of a continuous evolution of the interface between sliding events and of sudden degradation during largest events. This results in the progressive creation of a third-body (i.e. granular gouge) layer separating the two walls of the fault.

This view is confirmed by damage profiles across the fault, as provided in Fig. 8 (lower part). We observe that the central part of the fault is constituted of completely damaged material. Past this gouge layer, damage progressively decreases from complete degradation (damage=1) to almost intact rock (damage close to 0.2) following a decreasing exponential pattern. The influence of the confining pressure is highlighted. For values of $\sigma_3$ increasing from 45 to 180 MPa, the thickness of the gouge layer evolves from $\approx 40$ µm to $\approx 70$ µm, respectively. The decrease of damage with the distance from the fault is also faster for low confining stresses than for larger ones.
6 Perspectives

The model described in this study is able to reproduce typical seismic cycles (frequency and amplitude of events) only relying on independent calibrations of its parameters, and so without using any ad-hoc slip-weakening or rate-and-state friction law at contacts between elements composing the fault. Upon further use and interpretations, it is thus expected to provide valuable information on the local phenomena at stake during laboratory earthquakes.

Future work will be dedicated to overcome the current limitations of the model by increasing its size, by introducing a realistic roughness at the fault surface, and by implementing plasticity in the bulk rock continuum model, observed in [1, 2]. More data about stress redistributions and fault damage evolution, as well as detailed statistics on acoustic and slip events, will also be gathered in order to analyse more deeply the influence of gouge emission on the fault rheology.

References

[1] J. Aubry, F.X. Passelègue, J. Escartín, J. Gasc, D. Deldicque, A. Schubnel, J. Geophys. Res.: Solid Earth 125, e2020JB019670 (2020)
[2] J. Aubry, Séismes au laboratoire: friction, plasticité et bilan énergétique, PhD Thesis, ENS, PSL (2019)
[3] G. Mollon, Tribology Inter. 90, 343-355 (2015)
[4] J.T. Fredrich, B. Evans, T.-F. Wong, J. Geophys. Res.: Solid Earth 94, 4129-4145 (1989)
[5] G. Mollon, Comput. Part. Mech. 5, 517–527 (2018)