Characterization of variable-sensitivity force sensor using stiffness change of shape-memory polymer based on temperature

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Abstract
In the present study, we propose a variable-sensitivity force sensor using a shape-memory polymer (SMP), the stiffness of which varies according to the temperature. Since the measurement range and sensitivity can be changed, it is not necessary to replace the force sensor to match the measurement target. Shape-memory polymers are often described as two-phase structures comprising a lower-temperature “glassy” hard phase and a higher-temperature “rubbery” soft phase. The relationship between the applied force and the deformation of the SMP changes depending on the temperature. The proposed sensor consists of strain gauges bonded to an SMP bending beam and senses the applied force by measuring the strain. Therefore, the force measurement range and the sensitivity can be changed according to the temperature. In our previous study, we found that a sensor with one strain gauge and a steel plate had a small error and a large sensitivity range. Therefore, in the present study, we miniaturize this type of sensor. Moreover, in order to describe the viscoelastic behavior more accurately, we propose a transfer function using a generalized Maxwell model. We verify the proposed model experimentally and estimated the parameters by system identification. In addition, we realize miniaturization of the sensor and achieve the same performance as in our previous study. It is shown that the proposed transfer function can capture the viscoelastic behavior of the proposed SMP sensor quite well.

Keywords: Shape-memory polymer, Force sensor, Glass transition temperature, Cantilever, Strain gauge, System identification, Viscoelasticity

Introduction
Force sensors have been applied to various fields and are required to measure wider load ranges. One example in industry is a manufacturing system that has the flexibility to cope with various kinds of small-quantity production referred to as a flexible manufacturing system. Moreover, in rapidly aging societies, robotic technology has been applied to various fields, including industrial fields as well as nursing and welfare fields [1]. In these applications, wide-range force sensors that can obtain load information can measure multiple biosignals, such as heart rate, respiration cycle, and weight transitions [2]. Most force sensors transform the mechanical deformation of the detection area under an applied force into a change in resistance, capacitance, or reflectance that can be measured using electric signals. For example, some force sensors consist of strain gauges bonded to a bending beam. However, with this approach, it is difficult to change the measurement range or sensitivity of a sensor, both of which depend on the material used, the type of strain gauge, and the measurement method. The deformation range depends on the sensor material, and it is difficult to change these specifications after the sensor is produced. For this reason, we previously developed a force sensor using a
Shape-memory polymer (SMP), the measurement range and sensitivity of which can be changed [3, 4].

Shape-memory polymers [3–19] are increasingly being investigated as smart materials and are used in various fields. Shape-memory polymers change their modulus around the glass transition temperature \( T_g \), and are often described as two-phase structures comprising a lower-temperature “glassy” hard phase and a higher-temperature “rubbery” soft phase. The hard and soft phases are characterized by two different elastic moduli: an elastic modulus for the lower-temperature, higher-stiffness glassy plateau and an elastic modulus for the higher-temperature, lower-stiffness rubbery plateau. The reversible change in the elastic modulus between the glassy and rubbery states of SMPs can be as high as several 100-fold.

Since the stiffness of the SMP can be changed according to the temperature, the measurable force range determined based on the above strain range can also be changed. Moreover, even if the strain resolution is the same, the force resolution can be changed in a similar manner. In this way, the measurement range and sensitivity of the force sensor can be changed according to the temperature. Generally, environmental stability is an important requirement for sensors. For example, special compensating elements are often incorporated either directly into sensors or into signal conditioning circuits in order to compensate for temperature errors [20]. Note that, inverting the above idea, the proposed sensor uses the temperature-dependent changes positively.

In previous studies [3, 4], we made several prototypes of this sensor by attaching a strain gauge to an SMP sheet with an embedded electrical heating wire and evaluated their basic characteristics. Through experiments with these prototypes, which use the stiffness change of the SMP based on the temperature, we showed that the measurement range and sensitivity can be changed without replacing the actual sensor [3]. On the other hand, the changes in measurement range and sensitivity (ranging from a 100-fold to a thousand-fold) depend on the Young’s modulus of the SMP and are not adjustable. However, the change may be too large for some applications. Therefore, we also affixed a thin steel plate (see “Basic concept of force sensor” section for additional details). Furthermore, we could reduce the influence of the difference in elastic modulus between the strain gauges and the SMP sheet. This made it possible to reduce the discrepancy between the theoretical values and the measured values. Moreover, SMP force sensors with either one or two strain gauges and steel plates were fabricated, and their accuracy and sensitivity were investigated under the same conditions [4]. Experiments using the prototypes demonstrated that a sensor with one steel plate had a small error and a large sensitivity range, although the dimensions of the sensors were not optimized.

Therefore, in the present study, we miniaturized this type of sensor for practical use. A prototype of this sensor was made by attaching a strain gauge to an SMP sheet with an embedded electrical heating wire, and we evaluated the basic characteristics of the prototype sensor.

In a previous study, we proposed a theoretical formula that took the viscosity of the SMP into consideration, which made it possible to reduce the effect of stress relaxation [4]. However, since the first derivative of the measurement values was used to estimate the force (see “Basic concept of force sensor” and “Generalized Maxwell model” sections for additional details), the measurement errors could be large. Therefore, in the present study, we propose a transfer function method using a generalized Maxwell model. We verified the proposed model experimentally and estimated the parameters by system identification.

**Basic concept of force sensor**

As shown in Fig. 1, the force sensor was fabricated by bonding an SMP sheet and a steel plate and attaching a strain gauge to the steel. If the viscosity of the SMP and the embedded wire are negligible, by assuming a composite beam consisting of SMP and steel sheets, then the strain on the strain gauge \( \varepsilon \) can be expressed as follows [3, 4, 21]:

\[
\varepsilon = \frac{E_p h_p x}{I_p h_s^2} - \frac{E_s h_s x}{I_s h_s^2}
\]

![Fig. 1 Bending state of beam](image-url)
\[ W = \left( \frac{E_p l_p' + E_s l_s'}{h_p + h_s - h_1} \right) x \] (1)

where \( x \) is the distance between the strain gauge and the position at which the force is applied, \( E_p \) and \( E_s \) are the elastic moduli of the SMP and steel, respectively, and \( I_p' \) and \( I_s' \) are the area moments of inertia for the SMP and steel, respectively, and the neutral axis of the composite beam. Here, \( I_p' \) and \( I_s' \) are expressed as follows:

\[ I_p' = \frac{bh_p^3}{12} + \left( h_1 - \frac{h_p}{2} \right)^2 bh_p \] (2)

\[ I_s' = \frac{bh_s^3}{12} + \left( h_p - h_1 + \frac{h_s}{2} \right)^2 bh_s \] (3)

where \( b \) is the width of the beam, \( h_p \) and \( h_s \) are the thicknesses of the SMP and steel plates, respectively, and \( h_1 \) is the distance between the neutral axis and the SMP surface and is expressed as follows:

\[ h_1 = \frac{E_p h_p^2 + E_s (2h_p h_s + h_s^2)}{2(E_p h_p + E_s h_s)} \] (4)

As described in “Introduction” section, \( E_p \) can be changed according to the temperature. Therefore, the relationship between \( \varepsilon \) and \( W \) given in Eq. (1) can also be changed. Therefore, the change in the measurement range based on temperature can be modified. Moreover, as shown in Eq. (1), by changing the thickness of the steel plate, the measurement range and the sensitivity of the sensor can be modified.

In order to understand the mechanisms producing such unique properties and to design products including SMPs, various mathematical models have been proposed [5, 6, 12–19, 22]. To investigate relaxation processes in polymers, combinations of elements, including springs and dashpots, are widely used for modeling under isothermal conditions [12–18, 22]. For example, Tobushi et al. proposed a linear constitutive model by modifying a three-element viscoelastic model combining two springs and a dashpot to represent the deformation characteristics of the SMP [12]. Similarly, in our previous study [4], considering the viscosity of the SMP, we assumed that the relationship between \( W \) and \( \varepsilon \) is given by

\[ \dot{W} = L \left( \frac{\dot{\varepsilon}}{M} + \frac{\varepsilon}{N} \right) \] (5)

where \( L, M, \) and \( N \) are constants.

### Table 1 Specifications of force sensor in previous and present studies

|                   | Previous study [4] | Present study |
|-------------------|--------------------|--------------|
| \( l \) (mm)      | 63                 | 14           |
| \( b \) (mm)      | 20                 | 8            |
| \( h_p \) (mm)    | 1                  | 1.5          |
| \( h_s \) (mm)    | 0.05               | 0.07         |

### Table 2 Characteristics of SMP MP4510

| Properties                  | MP4510 |
|-----------------------------|--------|
| Elongation (\( T < T_g \))  | \( \geq 10\% \) |
| Elongation (\( T > T_g \))  | \( \geq 5\% \) |
| Elongation (\( T < T_g \))  | 10–30  |
| Elongation (\( T > T_g \))  | 25–30  |

### Miniaturization of sensor

#### Design of miniaturized sensor

The dimensions of the prototype SMP sensors used in our previous study [4] and the present study are shown in Table 1. Here, \( l \) is the length of the sensor. The dimensions of the prototype sensor constructed in the present study are similar to those of the commercial force sensor (LVS-2KA (\( T < T_g \)) and LVS-200GA (\( T > T_g \)), Kyowa Electronic Instruments Co., Ltd.) used in the experiment. The volume of the proposed sensor was reduced by 87% compared with that in our previous study. In our previous study, we applied a deformation of 5 mm to the tip of the sensor. In the present study, we determined \( b \) and \( h_p \) so that the reaction force \( W \) is the rated capacity of a commercial load cell (20 N) when the deformation is 1 mm below \( T_g \). The ratio of the sensor length (14 mm) to the deformation (1 mm) is approximately the same as that in the previous study, because the large creep strain of the SMP below \( T_g \) is not recovered [7]. The thickness of the steel plate \( (h_s = 0.07 \text{ mm}) \) is similar to that in our previous study [4]. The relationship between the applied force \( W \) and the deflection \( y \) is expressed as follows [21]:

\[ y = \frac{W l^3}{3 \left( E_p l_p' + E_s l_s' \right)} \] (6)

where \( l (= 11 \text{ mm}) \) is the distance between the fixed end and the position at which the force is applied. In the present study, we chose a polyurethane SMP (SMP Technologies Inc., MP4510, \( T_g = 45 \degree \text{C} \)).
fundamental characteristics of this material taken from the product catalogue are listed in Table 2. Substituting $E_p = 1,350 \text{ MPa ($T < T_g$)}$ and $E_s = 193 \text{ GPa}$ into this equation, we obtain $W = 27 \text{ N}$.

**Prototype**

In this study, we prepared two prototype sensors (samples i and ii) to check the reproducibility of the sensor fabrication process. The prototype SMP force sensor is shown in Fig. 2. In the present study, we prepared an SMP sheet with an embedded electrical heating wire in a manner similar to that described in our previous studies [4, 8]. The shape and dimensions of the heating wire (nichrome, outer diameter: 0.26 mm, electrical resistivity: $108 \times 10^{-6} \text{ Ω cm}$) are shown in Fig. 3. The underlined lengths in the figure were smaller than those in our previous study.

We bonded the SMP sheet and steel plate (SUS304H, thickness: 0.07 mm) using an adhesive (PPX, Cemedine Co., Ltd.). We attached one strain gauge (KFGS-2-120-C1-16 L1M2R, Kyowa Electronic Instruments Co., Ltd.) to the steel plate and measured the strain on the surface of the steel plate. The distance between the strain gauge position and the position at which the force was applied ($x$) was 6 mm. We used a cyanoacrylate adhesive (CC-36, Kyowa Electronic Instruments Co. Ltd., operating temperature range: −30 to 100 °C). The size and electrical resistivity of the prototype sensors are listed in Table 3. Since we fabricated the sensors manually, there were several manufacturing errors. In future studies, it

![Fig. 2 Prototype force sensor using an SMP sheet with an embedded electrical heating wire (sample i). We attached a thermocouple to the surface of the SMP sheet at the position indicated by the red circle](image1)

![Fig. 3 Heating wire. The underlined lengths are smaller than those in our previous study](image2)

**Table 3** Size of prototype force sensors

|                      | #i          | #ii         |
|----------------------|-------------|-------------|
| Length ($a$) (mm)    | $30.5 \pm 0.1$ | $29.9 \pm 0.1$ |
| Width ($b$) (mm)     | $8.5 \pm 0.0$  | $8.0 \pm 0.1$  |
| Thickness ($h_0 + h_f$) (mm) | $1.8 \pm 0.0$ | $1.9 \pm 0.0$ |
| Electrical resistivity ($\Omega$) | $1.8$ | $1.9$ |

* Average ± standard deviation from 10 measurements
will be necessary to minimize these errors by using a jig and a manufacturing machine.

A schematic diagram of the temperature control system is shown in Fig. 4a. The heating wire was connected to a power supply (PE18-1.3AT, KENWOOD) with a stable direct current, and a voltage of 5 V was supplied to the prototype sensor. With a thermocouple attached to the surface of the SMP sheet (red circle in Fig. 2), we heated the sheet and maintained a temperature of 70 °C (above \( T_g \)) as in our previous study [4]. Note that we used only two temperatures above and below \( T_g \) (70 °C and room temperature, respectively), and maintained these temperatures in order to minimize creep and thermal expansion. A pulse width modulation (PWM) signal was controlled by software developed in LabVIEW (National Instruments Co.). When the SMP was heated from room temperature to 70 °C, we set the duty ratio to 100% and then modified it to compensate for the heat loss and maintain the SMP temperature at 70 °C. The changes in temperature and strain with time for samples i and ii are shown in Fig. 4b, c, respectively. After keeping the samples at room temperature for 100 s, the temperature was increased and maintained at 70 °C by PWM control. The strain below and above 70 °C was almost constant although the values were very different due to thermal expansion of the SMP sheet. The difference in the strain above \( T_g \) for the two sensors is attributed to the manufacturing errors shown in Table 3. The fluctuation of the strain at 70 °C is thought to be due to the viscoelasticity of the SMP sheet, which caused a change in its shape. In the experiments described in the “Experiments” section, we evaluated the prototypes after the temperature became almost constant.

We captured a thermogram of the heated SMP sheet shown in Fig. 5a using an infrared thermal camera (NEC Avio Infrared Technologies Co., Ltd., F30W), and the results are shown in Fig. 5b. The entire sheet was heated uniformly to approximately 70 °C. One reason for the temperature uniformity is the reduced distance between the heating wire segments. Since the tape covered the surface to attach the thermocouple, the center of the SMP sheet is shown in blue.

Step 1: The sensor was held motionless in the unloaded state (just before touching).

Step 2: After the unloaded state, the sensor was moved in the direction of the blue arrow in Fig. 7 and was brought into contact with the load cell to apply a deformation of 1 mm to the tip of the sensor.

Step 3: Leaving the tip deformed, the sensor was held motionless.

Step 4: The sensor was returned to the initial position.

Steps 1 through 4 were repeated. The number of repetitions was 5, which was larger than in our previous study [4]. For system identification, the input should be “persistently exciting”, i.e., it should contain many distinct frequencies [23]. Therefore, in the present study, considering the potential applications of our sensor (i.e., a wide range of inputs), we randomly set the velocity in Steps 2 and 4 from 0.5 to 5 mm/s, and the rest time in Steps 1 and 3 from 0 to 20 s. For each condition, the measurements were conducted six times. We evaluated two prototype
Fig. 4  
(a) Schematic diagram of temperature control system. Change in surface temperature of SMP force sensor (b #i and c #ii)
sensors (samples 1 and 2). We prepared samples 1 and 2 by attaching samples i and ii, respectively, between the presser and base plates, as shown in Fig. 6b.
Experiment 2

We then evaluated the dynamic response of the proposed sensor to a step deformation. Similarly to experiment 1, the sensor was deformed as follows:

Step 1: The sensor was held motionless in the unloaded state (just before touching) for 10 s.
Step 2: After the unloaded state, the sensor was moved in the direction of the blue arrow in Fig. 7 (5 mm/s) and was brought into contact with the load cell in order to apply a deformation to the tip of the sensor.
Step 3: Leaving the tip deformed, the sensor was held motionless for 300 s.

We set the deformation in Step 2 to 0.25, 0.5, 0.75, or 1 mm in order to check whether the force increases with increasing deformation. For each condition, the measurements were conducted three times. We evaluated two prototype sensors (samples 1’ and 2’). After experiment 1, we removed the sensors from the presser and base plates, and reassembled samples 1’ and 2’ manually, as shown in Fig. 6b. Therefore, I in Fig. 2 was not exactly the same for samples 1 (2) and 1’ (2’).

Transfer function

Generalized Maxwell model

Here, \( \dot{\varepsilon} \) and \( \varepsilon \) are measured by the strain gauges, so \( \dot{W} \) is determined using Eq. (5) and the value of \( \dot{W} \) at each time step. If \( W=0 \) at \( t=0 \), the value of \( \dot{W} \) determined from Eq. (5) can be substituted into Eq. (7) in order to sequentially calculate \( W \):

\[
W(t + dt) = \dot{W}(t) + W(t) dt \tag{7}
\]

In our previous study [4], after performing this procedure multiple times, the force measured using the load cell was compared with \( W \) obtained based on the strain measured using the strain gauge, and the error was calculated. Then, \( W \) was determined by combining the theoretical equations in “Basic concept of force sensor” section. In MATLAB, the least-squares method was then used to determine the optimum values of \( L, M, \) and \( N \) in Eq. (5).

In the present study, in order to describe the viscoelastic behavior more accurately, we derived a transfer function using the generalized Maxwell model, as shown in Fig. 8 [13]. Westbrook developed a generalized Maxwell model to capture the shape-memory effect using two sets of nonequilibrium branches for two fundamentally different modes of relaxation: the glassy mode and the Rouse modes [13]. Since the temperature of our sensor is fixed above and below \( T_g \), we neglected the thermal expansion of the SMP. Then, using the transfer function, we calculated the force from the strain.

The derivation process is shown below. As shown in Fig. 8, the forces applied to each element \( f_0, f_i (i = 1 \cdots n) \) are expressed as follows:

\[
f_i = \frac{A E_i}{L_2} u_{is} = \frac{A\eta_i}{L_1} \dot{u}_{id} \tag{8}
\]

\[
f_0 = \frac{A E_0}{L} u \tag{9}
\]

where \( A \) is the cross-sectional area, \( u_{id} \) and \( u_{is} \) are the displacements of the dashpot and the spring, respectively, \( \eta_i \) is the viscosity, \( E_0 \) and \( E_i \) are elastic moduli, \( L \) is the total length of the proposed model, and \( L_1 \) and \( L_2 \) are the length of the dashpot and the spring, respectively. The total applied force \( f \) and the displacement \( u \) of the model are expressed as follows:

\[
f = f_0 + \sum_{i=1}^{n} f_i \tag{10}
\]

\[
u = u_{is} + u_{id} \tag{11}
\]

Using Eqs. (8) and (11), we have

\[
\dot{u} = \dot{u}_{is} + \dot{u}_{id} = \frac{L_2}{A E_i} \ddot{f}_i + \frac{L_1}{\eta_i} \dot{f}_i \tag{12}
\]

By Laplace transformation of Eqs. (9) and (12), we obtain

\[
F_0(s) = \frac{A E_0}{L} U(s) \tag{13}
\]
Fig. 9 Forces measured using prototype sensor and load cell when random force was applied (sample 1). FIT values are shown in each figure. a $T < T_g (n = 1)$. b $T > T_g (n = 1)$. c $T < T_g (n = 2)$. d $T > T_g (n = 2)$. e $T < T_g (n = 3)$. f $T > T_g (n = 3)$
\[ sU(s) = (B_i s + C_i)F_i(s) \] (14)

where

\[ B_i = \frac{L_i}{AE_i} \] (15)

\[ C_i = \frac{L_i}{AE_i} \] (16)

After Laplace transformation of Eq. (10), by substituting Eqs. (13) and (14), we have

\[ F(s) = F_0(s) + \sum_{i=1}^{n} F_i(s) \]

\[ = U(s) \left( \frac{AE_0}{L} + \sum_{i=1}^{n} \frac{s}{B_i + C_i} \right) \] (17)

Using Eq. (17), the transfer function \( G_i(s) \) for the generalized Maxwell model, shown in Fig. 8, with an input \( U(s) \) and an output \( F(s) \), is expressed as follows:

\[ G_i(s) = \frac{AE_0(B_i + C_i)(B_2s + C_2) \cdots (B_5s + C_5)}{L_i(B_2s^2 + C_2s + 1) \cdots (B_5s + C_5)} + \frac{+sL_i}{(B_2s^2 + C_2s + 1) \cdots (B_5s + C_5)} \] (18)

As shown in Eq. (18), the numbers of poles and zeros are \( n \). In the present study, we assumed that the transfer function \( G(s) \) with an input \( e(s) \) and an output \( W(s) \) had a structure similar to \( G_i(s) \) and \( n \) poles and \( n \) zeros. Namely, \( G(s) \) is expressed as follows:

\[ G(s) = \frac{s^{n} + \alpha_1 s^{n-1} + \cdots + \alpha_n}{s^{n} + \beta_1 s^{n-1} + \cdots + \beta_n} \] (19)

Since the temperature of the proposed sensor is fixed above or below \( T_g \), we assumed that \( \alpha_i (i=0 \cdots n) \) and \( \beta_i \) \((i=0 \cdots n)\) are constants in the rubbery and glassy states. Note that by Laplace transformation of Eq. (5), \( G(s) \) with \( n=1 \) can be obtained. We determined \( n \) using the experimental results.

**Calculation of force**

Using the results of experiments 1 and 2 shown in “Experiments” section, we first estimated the transfer function shown in Eq. (19) and calculated the fit ratio [24]. The fit ratio (hereinafter denoted as FIT) is defined as

\[ \text{FIT} = \left( 1 - \frac{\sum_{k=1}^{m} (\hat{y}(k) - y(k))^2}{\sum_{k=1}^{m} (\hat{y}(k) - y)^2} \right) \times 100\% \] (20)

where \( y(k) \) and \( \hat{y}(k) \) are the measured output and the simulated output, respectively, at time \( k \), \( \bar{y} \) is the average value of \( y \), and \( m \) is the number of samples. When \( y(k) \) is identical to \( \hat{y}(k) \) for all \( k=1, \ldots, m \), FIT becomes 100. The procedure is as follows:

- Method 1 (proposed in the present study): We estimated \( G(s) \) \((n=1–5)\) using the tfest() function in MATLAB. For the estimation, we imposed the condition that \( \alpha_i \) and \( \beta_i \) are positive because the values of \( L, A, E_0, B_i, \) and \( C_i \) in Eq. (18) are positive. The initial values of \( \alpha_i \) and \( \beta_i \) were set to 1 or NaN. The use of NaN indicates unknown coefficients. First, regarding each experimental data, we determined \( \alpha_i \) and \( \beta_i \) with maximum FIT value as the optima. Second, we calculated the average values of \( \alpha_i \) and \( \beta_i \) below and above \( T_g \) and determined \( G(s) \) with \( n=1–5 \). We then calculated FIT using the compare() function to validate the
determined \( G(s) \). Comparing FIT for each condition, we determined the optimal value of \( n \).

Furthermore, using the method in our previous study [4] (namely Eqs. (5) and (7), method 2) and the results of experiment 1, we estimated \( W \), calculated FIT, and compared the obtained results with those measured by Method 1. The procedure is as follows.

- Method 2 (used in our previous study): Using Eqs. (5) and (7), we calculated \( L \), \( M \), and \( N \). In MATLAB, the least-squares method was used to determine the optimum values of \( L \), \( M \), and \( N \) in Eq. (5). Here, \( L \), \( M \), and \( N \) were changed at intervals of 10 N s\(^{-1} \), 10 Ns, and 0.1 s, respectively. Based on these values, we also calculated FIT.

**Results and discussion**

**Identification of transfer function (experiment 1, method 1)**

Typical transitions of the force below and above \( T_g \) are shown in Fig. 9. As shown in Fig. 9, the two measured forces are almost identical. Moreover, by considering the viscosity of the SMP, the estimated force can reproduce the stress relaxation phenomenon. The maximum force and strain for different conditions are shown in Fig. 10. Although we applied a similar displacement below and above \( T_g \), the measured force range is significantly different. Based on the above results, it was shown that the miniaturized sensor achieved the same performance as in our previous study. On the other hand, the maximum strain below and above \( T_g \) is similar. The measured force below \( T_g \) is smaller than expected (27 N, see “Design of miniaturized sensor” section for additional details) and both the
maximum force and strain are different for the two prototype sensors. These differences may be attributed to errors during sensor manufacture, as shown in Table 3.

On the other hand, when the sensor was returned to the initial position, the force measured by the load cell became zero although the estimated values were not zero (arrow in Fig. 9b). One reason is that the SMP sheet could not recover to the initial shape quickly because of its viscosity, and the indenter of the load cell could not contact the prototype sensor. Above $T_g$, a fluctuation of the estimated force can be seen. The transitions of the measured strain are shown in Fig. 11. Similarly to Fig. 9, fluctuations can also be seen, and are attributed to electrical noise in the heating wire.

Using Method 1 ($n=1$–5), we calculated FIT. The mean ± the standard deviation below and above $T_g$ are shown in Fig. 12a, b, respectively. When $n=4$ and 5, by imposing the condition that the initial values of $a_i$ and $b_i$ are both 1, an error occurred during the calculation in MATLAB. When $n=5$, below $T_g$, there was a case when FIT was negative value, and the average value was small, as shown in Fig. 12a. The values of FIT below $T_g$ are larger than those above $T_g$. One reason would be the electrical noise shown in Figs. 9 and 11. As shown in Fig. 12, except for $n=5$, as the order of the transfer function became large ($n$ increased), the FIT values also increased.

We then calculated the average values of the coefficients in Eq. (19) below and above $T_g$, and determined
$G(s)$ with $n = 1–5$. Using the same experimental results, we calculated the mean and the standard deviation of FIT below and above $T_g$ for each $n$. The calculated FIT values are shown in Fig. 13; for most conditions, the largest value was obtained for $n = 3$. Therefore, a transfer function with $n = 3$ is considered to be optimal.

Bode diagrams for the transfer function ($n = 3$) obtained using Method 1 below and above $T_g$ are shown in Fig. 14. It can be seen that the gain has a significant temperature dependence. The phases above $T_g$ were larger than those below $T_g$, which can be attributed to the viscosity of the SMP. There were several differences between the gains and the phases for the two sensors, which may be attributed to the manufacturing errors shown in Table 3.

Comparison with our previous studies (experiment 1, comparison between Methods 1 and 2)
The FIT values determined using Method 2 are shown in Fig. 15. As shown in Figs. 12 and 15, the maximum values of FIT determined using Method 1 were larger than those determined using Method 2. Using the proposed transfer function model, the FIT values became larger than those in our previous studies.

Bode diagrams of the transfer function obtained using Method 2 below and above $T_g$ are shown in Fig. 16. The gains in Fig. 16 are similar to those in Fig. 14. However, above $T_g$, the phases in Fig. 16 are smaller than those in Fig. 14. This difference would cause a decrease in FIT.

Step deformation response (Experiment 2)
Using Method 1 ($n = 1–5$) and the results of experiment 2, we calculated FIT for different conditions. The mean ± the standard deviation below and above $T_g$ are shown in Fig. 17a, b, respectively. When $n = 4$ and 5, by imposing the condition that the initial values of $a_i$ and $b_i$ are both 1, an error occurred during the calculation in MATLAB. When $n = 4$, there were two cases when FIT was negative value, and its average value was small, as shown in Fig. 17. Similarly to experiment 1, Fig. 17 shows excellent agreement between the model estimations and the experimental data, although the deformations of the SMP sensors were not the same. As shown in Figs. 12 and 17, the values of FIT in experiment 2 were smaller than those in experiment 1. One reason would be the changes in the temperature and the sensor characteristics over time, as shown in Fig. 4b, c, because experiment 2 (more
than 310 s) was longer than experiment 1 (about 104 s). In the future, we will attempt to minimize these effects, for example by using improved temperature control, to determine the optimal value of \( n \).

We then calculated the average values of the coefficients in Eq. (19) below and above \( T_g \), and determined \( G(s) \) for \( n = 1 \sim 5 \). Using the same experimental results, we calculated the mean and the standard deviation of FIT below and above \( T_g \) for each \( n \). The calculated values of FIT are shown in Fig. 18, and are seen to be much lower than those in Fig. 17. The main reason is that the coefficients in Eq. (19) are different for each deformation. Therefore, for practical applications of the proposed sensor, it would be necessary to set the optimum coefficients according to the operating conditions.

**Conclusion**

We have developed a variable-sensitivity force sensor using an SMP sheet with an embedded electrical heating wire. In the present study, we miniaturized this type of sensor while referencing the dimensions and rated capacity of a commercial load cell. The volume was decreased by 87% compared with that in our previous study. The entire sheet of the prototype sensor was heated uniformly to approximately 70 °C.

Moreover, we proposed a transfer function using a generalized Maxwell model. Using identification experimental results, we determined the numbers of poles and zeros and compared the FIT value between our previous and present studies. Models were introduced and were validated experimentally, and there was excellent agreement between the model estimations and the experimental data. A transfer function with \( n = 3 \) was found to be optimal. Using the proposed model, the FIT value became larger than in our previous studies.

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Not applicable.
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