The Impact of Stellar Model Spectra in Disk Detection

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ABSTRACT

We present a study of the impact of different model groups in the detection of circumstellar debris disks. Almost all previous studies in this field have used KURUCZ (ATLAS9) model spectra to predict the stellar contribution to the flux at the wavelength of observation thus determining the existence of a disk excess. Only recently have other model groups or families like MARCS and NEXTGEN (PHOENIX) become available to the same extent as ATLAS9. This study aims to determine whether the predicted stellar flux of a disk target can change with the choice of model family - can a disk excess be present in the use of one model family whilst being absent from another? A simple comparison of KURUCZ model spectra with MARCS and NEXTGEN model spectra of identical stellar parameters was conducted and differences were present at near-infrared wavelengths. Model spectra often do not extend in wavelength to that of observation and therefore extrapolation of the spectrum is required. In extrapolation of model spectra to the Spitzer MIPS passbands, prediction of the stellar contribution differed by 5% at 70 µm for F, G and early K spectral types with differences increasing to 15% for early M dwarfs. Analysis of the Spitzer MIPS 24 µm observations of 37 F, G and K solar-like stars in the Pleiades cluster was conducted. In using KURUCZ model spectra, 7 disk excesses were detected while only 3 and 4 excesses were detected in using MARCS and NEXTGEN (PHOENIX) model spectra respectively.

1 INTRODUCTION

The study of Protoplanetary and Debris Disks is a significant area of research today. It is uncertain whether the processes which formed our own solar system are typical of other planetary systems. Both protoplanetary and debris disks describe planetary systems in the midst of formation/evolution and therefore are crucial to our understanding of planetary systems as a whole. In steady state, the mutual collisions of planetesimals replenish dust removed from the disk by radiation pressure and Poynting-Robertson drag. Planetesimals and regenerated dust around older stars describe debris disks in contrast to the first-generation dust around stars younger than 10 Myr old.

Detection of a disk may be gained if infrared emission from an object is in excess of expectation from the stellar flux. This is determined numerically by fitting a model spectrum and extrapolating out to the wavelength of observation - see the Appendix of Bryden et al. (2006) for details of this procedure. The excess ratio \( \frac{F_{\text{obs}}}{F_{\text{pred}}} \) is computed and detection of an infrared excess is concluded when this value exceeds the 3σ limit of the entire sample (Aumann et al. (1984), Beichman et al. (2005)). In almost all past and present studies of debris disks, KURUCZ (ATLAS9) models (Castelli & Kurucz (2004)) are used for this procedure. Other model groups like MARCS (Gustafsson et al. (2008)) and NEXTGEN (PHOENIX) (Hauschildt et al. (1999)) remain largely unused and have only recently become readily available.

Standard 1D model spectra are the result of a solution of the equation of radiative transfer, the hydrostatic equation, mixing-length theory and the equations describing the gas phase in chemical equilibrium. Ideally, every possible atomic and molecular transition should be included however this is computationally impractical and simplifications are necessary. MARCS and NEXTGEN model spectra are produced using Opacity Sampling (Helling & Jørgensen (1998)) where opacities are distributed statistically in wavelength giving rise to high resolution model spectra. KURUCZ model spectra however are produced using Opacity Distribution Functions (ODFs) where absorption cross-sections are binned producing a smoother, lower resolution model spectrum (Gustafsson et al. (1975), Ekberg et al. (1986)). Discrepancies in line features are expected between model groups due to these different opacity treatments though such differences are not expected at the longer wavelengths at which disks are observed. The different model groups include opacities drawn from different literature. This is likely to introduce sizeable differences in the spectral appearance of the model groups; for a discussion, see Jørgensen (2003). Model spectra with identical stellar parameters but of different model families are compared in Section 2 of this paper to determine whether such differences exist. The model spectra often do not extend to longer wavelengths at which observations of disks are
made therefore, fitting of a Rayleigh-Jeans tail and extrapolation to the required wavelength is necessary (Section 3). NextGen models do extend beyond 100 μm however, the sparse wavelength distribution at these wavelengths make it necessary to fit a Rayleigh-Jeans tail for interpolation onto some filter bandpass. Section 4 applies the use of all three model families to 37 solar-like stars previously analysed for disk excitation in Sierchio et al. (2010).

2 COMPARISON OF MODEL SPECTRA

Spectra of all three model families were compiled. KURUCZ and NEXTGEN model spectra with inclusion of a 2.0 km s⁻¹ microturbulence, a convective mixing length, l/H = 1.25 and no parameterization of convective overshoot were chosen (Castelli & Kurucz (2004), Hauschildt et al. (1999)). MARCS model spectra with identical microturbulence but (the only available at the time) convective mixing, l/H = 1.5 were chosen (Gustafsson et al. (2008)). KURUCZ and NEXTGEN model groups provide accurate synthetic SEDs of solar-like stars though KURUCZ better reproduces early-type stars while NEXTGEN is more accurate in the M regime (Bertone et al. (2004)). The Marcs model group have also been shown to adequately match the stellar SEDs given their use in the photometry calibration of the Spitzer Space Telescope (Decin et al. (2004)). Fig. 1 displays the effective temperature and surface gravity combinations of all three model families. As shown, model spectra in all three families adequately span the F, G and K spectral types which dominate the stellar sample of disk surveys. Model spectra of solar metallicity are considered the most appropriate in the context of disk surveys using stellar samples of solar-like, solar neighbourhood stars which typically exhibit metallicities of [Fe/H] = 0.25 (Valenti & Fischer (2005)).

A comparison of the model spectra across the model families was conducted. A KURUCZ model spectra of some effective temperature, surface gravity (and solar metallicity) was compared with a MARCS and NEXTGEN model spectra with identical stellar parameters. The wavelength range and resolution differed between model families therefore interpolation (Eq. 3) was necessary. In comparing KURUCZ and MARCS spectra, the former were interpolated onto the wavelength distribution of the latter. For comparison with NEXTGEN models, a few short wavelength points were omitted to allow a common wavelength range and the KURUCZ spectra interpolated onto the NEXTGEN wavelength distribution. A ratio of the flux (2) was computed.

\[
F_{i}' = F_{i-1} + (\lambda_j - \lambda_{i-1}) \frac{(F_i - F_{i-1})}{(\lambda_i - \lambda_{i-1})} \quad (1)
\]

\[
F_{\text{ratio}}(\lambda) = \frac{F_j}{F_i}' \quad (2)
\]

where \(\lambda_i, F_i\) represents the KURUCZ model spectrum while \(\lambda_j, F_j\) represents the MARCS or NEXTGEN model spectrum where \(\lambda_j\) is a wavelength intermediate to \(\lambda_{i-1}\) and \(\lambda_i\). \(F_{\text{ratio}}(\lambda)\) was calculated for each set of synthetic spectra and subsequently averaged for all spectra of the F, G and K spectral types.

Spectral fluxes of KURUCZ and MARCS spectra can differ by greater than an order of magnitude (Fig. 2). In particular, the poorest agreement occurs at approximately 0.2 μm. The trend of the flux ratio is dominated by delta-function-like features which likely are produced as a result of line features present in one model group but absent from the other. This is likely capturing the different opacity treatments. A higher level of agreement occurs between KURUCZ and NEXTGEN model spectra (Fig. 2). Line features remain present, again capturing the different opacity treatments. The trends of both comparison tend to unity with increasing wavelength showing better agreement of the synthetic spectra at longer wavelengths. The best agreement is apparent for the KURUCZ and NEXTGEN models while there exist some fluctuation about unity for the KURUCZ/MARCS comparison. The bump-like feature at approximately 10 μm in the KURUCZ/MARCS comparison possibly arises from differences in H⁻ opacity, the dominant source at these wavelengths. The deviation from unity in the KURUCZ/NextGen comparison arises from an unphysical anomaly present in the KURUCZ T = 8000 K, log(g) = 4.5, [M/H] = 0.0 model spectrum. This feature was masked in all subsequent analysis. These differences in the model spectra has obvious implications in stellar flux prediction and finding the excess ratio from which detection of a disk is decided.

3 EXTRAPOLATION TO LONGER WAVELENGTHS

The algebraic form of the Rayleigh-Jeans’ flux \((F_{RJ}(\lambda, T) = 2kTc/\lambda^4)\) did not provide a good fit of the long wavelength regions of the synthetic spectra. \(F_{RJ}(\lambda, T_{\text{eff}})\) (where \(T_{\text{eff}}\) is the effective temperature of the model spectrum) was offset

\[1\text{ http://kurucz.harvard.edu/grids.html}
\[2\text{ ftp://phoenix.hs.uni-hamburg.de/NextGen/Spectra/}
\[3\text{ http://marcs.astro.uu.se/search.php} \]
from the synthetic spectrum by up to 40% at 5 \( \mu m \). A simple scaling of \( F_{24, \mu m} \) to the synthetic spectrum did not offer a solution as the fall-off of the synthetic flux with wavelength deviated from \( \lambda^{-4} \). Such an offset and slope difference remained when the full blackbody flux was instead fit. A straight line fit of the spectrum in logarithmic space was instead conducted. Eq. 3 was fit (allowing freedom of \( a \) and \( b \) which minimised \( \chi^2 \)) and transformed back into normal space, giving a function of the form shown in Eq. 4:

\[
\log(F_\lambda) = a \log(\lambda) + b
\]

\[
F_\lambda = \frac{10^b}{\lambda^a}
\]

where \( a \) describes the wavelength power index (approximately 4) while \( b \) serves as a scale for the pattern. Significant differences in the parameters which minimised \( \chi^2 \) were apparent - for example, \( a \approx 3.9, b \approx 6.66 \) provides a good fit for a KURUCZ \( T = 3500 \) K, \( \log(g) = 4.5 \) and \([M/H] = 0.0 \) model spectrum while \( a \approx 4.03, b \approx 6.94 \) provides a good fit for a NEXTGEN model spectrum with the same parameters. For each model spectrum, such a fit was performed, the spectrum extrapolated out to 24 \( \mu m \) and 70 \( \mu m \) and the flux integrated over the corresponding Spitzer MIPS passbands which are sensitive to dust emission at 1 AU and outside 10 AU from a solar-like star respectively (Rieke et al. (2004)). For each set of KURUCZ and MARCS or NEXTGEN model spectra identical in stellar parameters, the broadband fluxes of these passbands were compared. Fig. 3 displays a ratio of these fluxes for each set of synthetic spectra. KURUCZ model spectra do not extend to the cooler effective temperatures of M dwarfs and therefore MARCS and NEXTGEN model spectra are also compared in this way to determine the possible differences in flux prediction for such objects.

At 24 \( \mu m \), the extrapolated fluxes of KURUCZ and MARCS model spectra exhibit good agreement (of approximately 2 \%) for F, G and early K spectral types. The trend deviates significantly from unity for cooler late K and early M spectral types with disagreement as high as 15 \% in KURUCZ and MARCS spectra of 3500 K. At 70 \( \mu m \), the flux ratio of KURUCZ and MARCS model spectra is identical to that of 24 \( \mu m \) though it has a larger deviation from unity (approximately 5\%). This suggests that the long wavelength regions of KURUCZ and MARCS synthetic spectra differ in slope and therefore their log-log straight line fits (Eq. 3) diverge with wavelength. In comparison, the KURUCZ/NEXTGEN trend exhibits a higher level of agreement which is consistent with Fig. 2. At 24 \( \mu m \), the flux ratio almost perfectly centres on unity with large disagreement only apparent in the comparison of early M dwarfs. At 70 \( \mu m \), a small deviation from unity is apparent though agreement is still true within 3 \% for F, G and K spectral types. Similarly, MARCS and NEXTGEN model spectra of F, G and K spectral types exhibit good agreement (of approximately 4 \%) at 24 \( \mu m \), deviating to approximately 8 \% at 70 \( \mu m \). Again, M dwarfs exhibit the worst agreement with differences as large as 20 \%. The poor agreement of M dwarf model spectra likely arises from the challenges in determining the opacity-relevant effects of molecular species present in such atmospheres (Gustafsson (1997), Jørgensen (2003), see also section 2.2 of Helling & Lucas (2009)). Observational studies have noted discrepancies between the observed and synthetic colours of such objects (Gautier et al. (2007)).

4 ANALYSIS OF 37 SOLAR-LIKE PLEIADES STARS

The possible impact in the use of different model groups in flux prediction of disk observations has so far only been inferred from comparison of the model spectra. It is therefore ideal to physically apply using all three model spectra to a set of observations.

A recent publication (Sierchio et al. (2010)) analysed the Spitzer MIPS 24 \( \mu m \) observations of approximately 70 stars in the Pleiades cluster. Using KURUCZ model spectra for prediction of the stellar flux at 24 \( \mu m \), they concluded that 23 of such stars harboured disks. Using a smaller subset of their target sample (due to limitations in the available photom-
Figure 3. A comparison of the Spitzer MIPS predicted fluxes against the effective temperature of the log(g) = 4.5 KURUCZ and MARCS/NextGen model spectra compared. The KURUCZ/MARCS comparison is shown as blue triangles, the KURUCZ/NextGen comparison are represented by red circles and the MARCS/NextGen comparison shown as the green squares. There are comparably fewer points for the KURUCZ and NextGen comparison as fewer spectra had matching stellar parameters. The dashed horizontal lines mark agreement of the fluxes.

Figure 4. A comparison of the predicted fluxes of the 37 targets (Sierchio et al. 2010) at 24 µm (in mJy) in using KURUCZ, MARCS and NextGen model spectra. Red circles show the comparison for KURUCZ and MARCS models while blue triangles show that of KURUCZ and NextGen models. The dashed line represents equality and is not a fit of either data set.
The distribution of excess ratios ($F_{\text{obs}} / F_{\text{pred}}$) at Spitzer MIPS 24 $\mu$m for the sample of 37 stars using KURUCZ (red), MARCS (blue) and NEXTGEN (green) model spectra as a means of predicting the stellar contribution. Gaussian fits are shown as solid black lines and the x-translation and deviation parameters of the fit ($\mu$ and $\sigma$) are also displayed. The $3\sigma$ confidence intervals associated with the Gaussians are shown as the vertical dashed lines. AKII437 and HII132 (with very significant excess ratios) are omitted for the purpose of sensible axes scales.

Detection of an infrared excess is considered conclusive when the excess ratio of an object exceeds the upper $3\sigma$ confidence interval of the entire distribution. A narrow distribution with a translation of $\mu = 1.0$ describes model spectra which accurately predict the flux of the target stars. While the MARCS distribution is best-centred on unity, it is the broadest of the three distributions. KURUCZ model spectra form the narrowest distribution while the NEXTGEN distribution has the largest offset from unity. These offsets have been previously noted of Pleiades stars - see Fig. 4 of Gorlova et al. (2006).

Figure 5. The distribution of excess ratios ($F_{\text{obs}} / F_{\text{pred}}$) at Spitzer MIPS 24 $\mu$m for the sample of 37 stars using KURUCZ (red), MARCS (blue) and NEXTGEN (green) model spectra as a means of predicting the stellar contribution. Gaussian fits are shown as solid black lines and the x-translation and deviation parameters of the fit ($\mu$ and $\sigma$) are also displayed. The $3\sigma$ confidence intervals associated with the Gaussians are shown as the vertical dashed lines. AKII437 and HII132 (with very significant excess ratios) are omitted for the purpose of sensible axes scales.

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5 SUMMARY

Detection of a disk can be concluded when infrared flux is in excess of expectation from the star alone. Popularly, a model spectra is used to predict the stellar contribution and almost all previous studies in this field have used KURUCZ (ATLAS9) model spectra for this purpose. MARCS and NEXTGEN model spectra however have remained largely unused due to only recent availability.

An initial comparison of solar metallicity KURUCZ, MARCS and NEXTGEN model spectra of identical stellar parameters yielded differences as large as two orders of magnitude. In extrapolating these spectra out to longer wavelengths, prediction of Spitzer MIPS fluxes could differ by 5% for F, G and K spectra types and as much as 15% for early M dwarfs.

A recent publication (Sierchio et al. (2010)) studied the Spitzer MIPS 24 $\mu$m observations of 70 stars in the Pleiades cluster for possible disk excesses, using KURUCZ (ATLAS9) models for prediction of the stellar contribution. This analysis was initially reproduced for a subset of these stars and 7 out of 37 were determined to have disk excesses. In repeating this analysis using MARCS and NEXTGEN model spectra, only 3 and 4 disk excesses were detected respectively. The differences between the model groups are such that different disk detection rates are obtained. We encourage the use of more than one model family to determine the presence (or absence) of an infrared excess.

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