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Key Points:
- Basic vertical profile measurements are sufficient to constrain entrainment in nonprecipitating Southern Ocean (SO) stratocumulus clouds
- The sensitivity of shortwave cloud forcing to entrainment is 2–20 times the sensitivity to particle concentration or cloud base updraft
- Quantifying the impact of entrainment on shortwave cloud forcing is needed to determine the energy budget in the SO

Supporting Information:
- Supporting Information S1

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Measured Constraints on Cloud Top Entrainment to Reduce Uncertainty of Nonprecipitating Stratocumulus Shortwave Radiative Forcing in the Southern Ocean

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Abstract Stratocumulus cloud top entrainment has a significant effect on cloud properties, but there are few observations quantifying its impact. Using explicit 0-D parcel model simulations, initialized with below-cloud in situ measurements, and validated with in situ measurements of cloud properties, the shortwave cloud radiative forcing (SWCF) was reduced by up to 100 W m−2 by cloud top entrainment in the Southern Ocean. The impact of entrainment-corrected SWCF is between 2 and 20 times that of changes in the aerosol particle concentration or updraft at cloud base. The variability in entrainment-corrected SWCF accounts for up to 50 W m−2 uncertainty in estimating cloud forcing. Measurements necessary for estimating the impact of entrainment on cloud properties can be constrained from existing airborne platforms and provide a first-order approximation for cloud radiative properties of nonprecipitating stratocumulus clouds. These measurement-derived estimates of entrainment can be used to validate and improve parameterizations of entrainment in Global Climate Models.

Plain Language Summary Clouds over the ocean have a significant impact on climate because they reflect sunlight that would otherwise be absorbed by the ocean. Understanding and accurately modeling how much sunlight these clouds reflect is important to understand feedbacks between clouds and climate. Using a simple model and cloud measurements, mixing of moist cloud air with warm-dry air from above the cloud was shown to decrease the cloud droplet number concentration and consequently total liquid water, which significantly decreases the overall amount of sunlight reflected by the cloud. Cloud droplets form onto small particles as they enter the base of the cloud through an updraft; in the meantime, warm-dry air from above the cloud mixes downward. In this study, the warm-dry air from above the cloud was shown to be more influential on reflecting sunlight than the concentration of particles and the updraft velocity at cloud base. These results emphasize the importance of accurately accounting for the mixing of warm-dry air from above the cloud in climate models and can be constrained by existing measurements that are readily available on weather balloons and aircraft.

1. Introduction

Low clouds in the Southern Ocean (SO) are poorly simulated by general circulation models (GCMs) and tend to overestimate the amount of radiation absorbed by the SO (Bodas-Salcedo et al., 2014; Haynes et al., 2011; Hyder et al., 2018; McCoy et al., 2014). This bias is driven by underestimates of cloud radiative forcing, which are likely due to errors in cloud microphysical properties such as droplet size and concentration, as well as cloud properties such as precipitation and entrainment (Mason et al., 2015; Vial et al., 2013). Cloud properties in the SO are also particularly sensitive to aerosol loading due to relatively low background aerosol concentrations (Downey et al., 1990; Whittlestone & Zahorowski, 1998). Therefore, addressing these uncertainties is necessary for assessing future climate change feedbacks. GCM shortcomings, biases in satellite retrievals, and the scarcity of SO cloud measurements (Lenschow et al., 1999; Russell et al., 1998; Seinfeld et al., 2016) motivated the Southern Ocean Clouds, Radiation, Aerosol Transport Experimental Study (SOCRATES), which conducted in situ airborne measurements of clouds, aerosol, and meteorological state over the SO on the NSF/NCAR HIAPER Gulfstream V (GV) (Laursen et al., 2006). The analysis shown here aims to
improve GCM estimates of reflected shortwave radiation and the Earth’s energy budget by obtaining measurements of cloud top entrainment drying with widely used measurements to improve the parameterization of clouds in models.

Aerosol-cloud interactions account for a significant part of the uncertainty in the global energy budget (IPCC, 2014; Lohmann, 2017; Nazarenko et al., 2017; Seinfeld et al., 2016; Stevens, 2015). Variability in reflected shortwave radiation (or albedo) is caused by differences in cloud fraction, cloud optical thickness, cloud droplet number concentration (CDNC), and cloud lifetime (Albrecht, 1989; Jiang et al., 2006; Wang et al., 2012). Aerosol indirect effects such as enhancements in particle concentrations leading to increased CDNC and increased reflected solar radiation are classic examples of how changes in particle concentrations alter cloud optical properties (Ackerman et al., 2000; Albrecht, 1989; Jiang et al., 2006; Lu & Seinfeld, 2005; Twomey, 1977; Xue & Feingold, 2006). While correctly representing CDNC is a concern in GCMs for accurately modeling the optical thickness of clouds, uncertainty in liquid water path (LWP) accounts for more than a factor of 2 change in cloud optical thickness compared to CDNC (Brenguier et al., 2011), making LWP of first-order importance (Boers & Ross, 1994; Lu et al., 2008). One of the biggest factors in determining marine stratocumulus LWP is the cloud top entrainment rate; however, numerous feedbacks between the cloud properties and entrainment make simulating and parameterizing entrainment rates challenging (Chen et al., 2011). Even though cloud top entrainment and detrainment play an important role in the resulting cloud optical properties, their roles in determining cloud optical properties are not often considered in aerosol-cloud closure studies.

Entrainment of free tropospheric air at cloud top typically results in a warmer, drier boundary layer, as well as clouds with subadiabatic LWP and decreased precipitation (Ackerman et al., 2009; Chen et al., 2011; Deardorff, 1980; Hill et al., 2009; Lu & Seinfeld, 2005; Wood, 2007; Xue & Feingold, 2006). Under certain thermodynamic conditions, when the above-cloud entrained air is excessively dry, the air mixture can result in a parcel becoming negatively buoyant due to evaporative cooling, leading to enhancements in the entrainment rate (Ackerman et al., 2004; Burnet & Brenguier, 2007; Grabowski, 1993). The exchange of air between cloud top and the overlying air is determined by the gradient in specific humidity and temperature across the inversion (Sanchez et al., 2017; Stevens, 2002), both of which are expected to change in a warming climate (Qu et al., 2015). Increased temperature gradients between the marine boundary layer (MBL) and free troposphere result in stronger inversions which decrease cloud top entrainment rates (Ackerman et al., 2004; Caldwell & Bretherton, 2009) and increase cloud fraction (Caldwell et al., 2013). However, enhanced specific humidity gradients across the inversion (Webb & Lock, 2013) also lead to increased evaporation-driven entrainment at cloud top, which thins stratocumulus clouds in LES simulations (Bretherton, 2015; Bretherton et al., 2013).

Cloud top entrainment is also affected by an aerosol-driven feedback on the entrainment rate. For example, a decrease in aerosol loading leads to larger cloud droplets and an enhancement in droplet sedimentation and precipitation at cloud top, which results in a decrease in cloud top LWC and, consequently, a decrease in evaporative cooling and entrainment as well (Bretherton, 2015; Bretherton et al., 2007; Turton & Nicholls, 1987; Zuidema et al., 2008). Large Eddy Simulation (LES) by Zuidema et al. (2008) shows increased particle concentrations are linked to decreased cloud fraction as a result of more entrainment in spite of less drizzle. Further complicating this process, cloud top entrainment is initiated by mixing on centimeter to meter length scales, meaning it cannot be explicitly resolved in LES models or GCMs, requiring entrainment to be a parameterized process. In previous LES studies of stratocumulus clouds, the vertical resolution has been increased in an attempt to reduce the uncertainty in the entrainment rate and resulted in a decrease in the overprediction of cloud top entrainment and, consequently, underpredicting the cloud radiative forcing (Bretherton et al., 1999; Stevens et al., 2005). In addition, Stevens et al. (2005) found that increasing the horizontal resolution is important for resolving the strength of large eddies (>70 m) and, hence, improves the calculation of the entrainment rate.

In GCMs, the Cloud Layers Unified by Binormals (CLUBB) parameterization is the most promising method of reproducing subgrid-scale cloud properties and involves using probability distribution functions (PDF) derived from subgrid-scale predictive moments to derive cloud properties over multiple regimes (Guo et al., 2010, 2014). GCMs are constrained with satellite measurements; however, current satellite-derived cloud properties assume a 30% reduction in CDNC due to inhomogeneous entrainment
In this manuscript, five case studies are used to quantify the impact of cloud top entrainment on stratocumulus cloud optical properties in the SO using a concept first described by Betts (1983) and Paluch (1979). Betts (1983) showed that stratocumulus clouds largely follow a mixing line structure from cloud base to cloud top, where conservative variables are linearly related, indicating that entrained air at cloud top vertically mixes throughout the cloud. This concept was recently utilized by Sanchez et al. (2017) and Calmer et al. (2019) to account for entrainment-induced evaporation of cloud droplets, which improved representation of nonprecipitating stratocumulus cloud optical properties. The method used in a pristine midlatitude marine environment (Sanchez et al., 2017) and a polluted Mediterranean region (Calmer et al., 2019) is also verified for SO clouds during the SOCRATES experiment. By combining in situ measurements of cloud condensation nuclei (CCN) spectra at cloud base, vertical profiles of the thermodynamic properties throughout the MBL, updraft, and an aerosol-cloud-parcel model (ACPM), this study obtains accurate simulations of the observed cloud optical properties by incorporating the impact of entrainment on the microphysical properties of stratocumulus clouds in the SO. We also relate the sensitivity of stratocumulus shortwave cloud forcing (SWCF) to changes in below-cloud aerosol concentrations and updraft and to the impact of entrainment ($ENT_{LWC}$; sections 2.4) via the reduction of LWP. The mixing-line entrainment approximation ($ENT_{ML}$; sections 2.5 and 3.2) enables the calculation of the entrained fraction of air and resulting subadiabatic LWC throughout the cloud using only vertical profiles of temperature, pressure, and relative humidity measurements. This technique to account for entrainment can eventually be applied globally to nonprecipitating stratocumulus clouds using relatively simple observations to constrain and evaluate GCM parameterizations.

2. Methods

2.1. NSF/NCAR HIAPER GV Measurements

Measurements used in SOCRATES are collected on the NSF/NCAR Gulfstream-V High-performance Instrumented Airborne Platform for Environmental Research (GV HIAPER) observational platform. The GV was stationed at the Hobart International Airport, Tasmania, during the austral summer between 15 January and 24 February 2018. The flight strategy during SOCRATES involved ferrying out to an area of interest followed by a series of straight vertical profiles, and level legs to sample below, in and above cloud in the MBL. The GV conducted 15 research flights (RFs) over the SO between 42.5°S and 62.1°S and between 133.8°E and 163.1°E at altitudes ranging from 50–7,500 m.

A wing-mounted Ultra-High Sensitivity Aerosol Spectrometer (UHSAS, DMT, Boulder, CO) measured particle size distribution between 0.06 and 1.0 μm in diameter; however, the 0.6–0.7 μm bins were not used due to instrument noise at these sizes. A condensation particle counter (CPC, TSI 3760A) was used to measure total particle concentrations ($N_p$, diameter >10 nm). CCN measurements were performed with two miniature continuous-flow streamwise thermal gradient chambers, one in scanning supersaturation mode and one in constant supersaturation mode (Roberts & Nenes, 2005). The supersaturation range in the scanning CCN counter spanned from 0.06–0.87%, and a single spectrum recorded every 5 min. The constant supersaturation CCN counter operated at 0.43% supersaturation, at 1 Hz. The internal chamber pressure of both CCN counters was controlled to 400 hPa. Updraft distributions at cloud base were measured by the GV HIAPER Cloud Radar (HCR), corrected for plane orientation with nadir looking data. HCR reflectance and a Particle Habit Imaging and Polar Scattering probe (Abdelmonem et al., 2016; Schnaiter et al., 2018) were also used to identify clouds without drizzle (Chin et al., 2000; Kato et al., 2001; Kogan et al., 2005). A high-frequency (25 Hz) Vertical Cavity Surface Emitting Laser (VCSEL) water vapor hydrometer was used to measure the water vapor mixing ratio (Zondlo et al., 2010). A QA/QC water vapor data set at 25 Hz for the SOCRATES campaign (Diao, 2020) is further used to derive specific humidity ($q_v$) and relative humidity with respect to liquid ($RH_{liq}$). Finally, a cloud droplet probe (CDP, DMT, Boulder, CO) is used to measure cloud droplet concentration and size (2–50 μm diameter) to calculate liquid water content (LWC) and cloud droplet extinction for comparison with the ACPM (section 2.2).
Table 1

GV Observed Below-Cloud Aerosol Concentrations, in-Cloud Droplet Properties, and Calculated Optical Properties

|                   | RF02  | RF10Thin | RF10Thick | RF12  | RF13  |
|-------------------|-------|----------|-----------|-------|-------|
| **Observed**      |       |          |           |       |       |
| CDNC_{obs} (cm^{-3}) | 64 ± 11 | 97 ± 17 | 94 ± 51 | 175 ± 34 | 143 ± 24 |
| SWCF_{obs} (W m^{-2}) | 103.1 | 142.9 | 340.1 | 234.3 | 201.1 |
| Droplet spectra width\(a\) | 0.81 ± 0.04 | 0.84 ± 0.03 | 0.73 ± 0.12 | 0.82 ± 0.05 | 0.83 ± 0.05 |
| Max LWC (g m^{-3}) | 0.23 | 0.30 | 0.68 | 0.40 | 0.36 |
| LWP (g m^{-2}) | 1.29 | 2.02 | 19.89 | 6.47 | 3.94 |
| Below cloud aerosol (cm^{-3}) | 457 ± 9 | 463 ± 62 | 956 ± 17 | 505 ± 52 | 290 ± 66 |
| Below cloud aerosol (>70 nm, cm^{-3}) | 161 ± 42 | 167 ± 29 | 193 ± 45 | 269 ± 74 | 199 ± 38 |
| Cloud base height (m) | 500 | 1,010 | 290 | 1,090 | 790 |
| Cloud top height (m) | 660 | 1,210 | 870 | 1,450 | 1,030 |
| Cloud base temperature (K) | 274.9 | 273.3 | 283.9 | 268.3 | 274.3 |
| Cloud top temperature (K) | 273.9 | 272.1 | 281.1 | 265.7 | 272.6 |
| **Simulations**    |       |          |           |       |       |
| CDNC_{adiabatic} (cm^{-3}) | 176   | 198     | 226      | 286   | 220   |
| SWCF_{adiabatic} (W m^{-2}) | 196.5 ± 17.1 | 220.8 ± 16.5 | 387.3 ± 9.4 | 272.7 ± 8.9 | 242.2 ± 7.0 |
| Droplet spectra width | 0.96 | 0.97 | 0.91 | 0.96 | 0.97 |
| Max LWC (g m^{-3}) | 0.32 | 0.38 | 1.59 | 0.55 | 0.47 |
| LWP (g m^{-2}) | 2.34 | 3.61 | 47.31 | 9.79 | 5.48 |
| TOA incoming solar radiation (W m^{-2}) | 453.8 | 419.4 | 431.3 | 248.1 | 210.8 |
| **LWC Entrainment (\(ENT_{LWC}\); Sanchez et al., 2017)\(e\) |       |          |           |       |       |
| CDNC_{ENTLWC} (cm^{-3}) | 79 ± 10 | 160 ± 5 | 166 ± 65 | 151 ± 20 | 132 ± 19 |
| SWCF_{ENTLWC} (W m^{-2}) | 135.2 ± 14.2 | 158.8 ± 15.2 | 345.4 ± 16.2 | 242.7 ± 10.4 | 212.5 ± 7.3 |
| 25th and 75th percentiles | 17.7–103 | 127–176 | 161–172 | 134–191 | 109–220 |
| CDNC_{ENTML} (cm^{-3}) 50th percentile | 79 ± 10 | 160 ± 5 | 166 ± 65 | 151 ± 20 | 132 ± 19 |
| 25th and 75th percentiles | 17.7–103 | 127–176 | 161–172 | 134–191 | 109–220 |
| SWCF_{ENTML} (W m^{-2}) | 135.2 ± 14.2 | 158.8 ± 15.2 | 345.4 ± 16.2 | 242.7 ± 10.4 | 212.5 ± 7.3 |
| 25th and 75th percentiles | 17.7–103 | 127–176 | 161–172 | 134–191 | 109–220 |
| **Difference between simulated and observed SWCF** |       |          |           |       |       |
| ΔSWCF_{adiabatic} | 93.4 ± 17.1 | 77.9 ± 16.5 | 47.2 ± 9.4 | 38.4 ± 8.9 | 41.1 ± 7.0 |
| ΔSWCF_{ENTLWC} | 32.1 ± 14.2 | 15.9 ± 15.2 | 5.3 ± 16.2 | 8.4 ± 10.4 | 11.4 ± 7.3 |
| ΔSWCF_{ENTML} | –0.9 ± 12.1 | 54.2 ± 16.2 | 27.6 ± 12.8 | –13.8 ± 11.0 | –11.8 ± 7.6 |

**Sensitivity Tests**

\( S(X_i) = \frac{d\ln(Y_i)}{d\ln(X_i)} \) (Feingold et al., 2003)

| \( X_i \) | \( Y_i = \text{SWCF} \) | \( Y_i = \text{CDNC}_{avg} \) |
|-----------|----------------|----------------|
| \( N_p \) | 0.149 (0.145) | 0.031 (0.059) |
| \( w \)   | 0.033 (0.034) | 0.014 (0.027) |
| LWP       | 0.632 | 0.132 |
|           | 0.776 | 0.707 |
|           | 0.043 | 0.485 |
|           | 1.014 | 0.776 |

**Note.** ACPM simulation results are present for adiabatic simulations, as well as simulations for LWC entrainment method and mixing line entrainment method. The sensitivity of below-cloud particle number concentration, updraft velocity, and cloud LWP (which is affected by cloud top entrainment) on SWCF and CDNC. Sensitivities for SWCF are calculated for both the adiabatic (\(ENT_{LWC}\)) simulations. Sensitivities for CDNC are the same for adiabatic simulations and \(ENT_{ML}\). A 28% and 58% reduction in LWP are used because these values represent the observed minimum and maximum percentage reduction in the adiabatic LWP, shown in the previous row.

\( a\)Martin et al. (1994). **Liquid water path is calculated from CDP distribution profile.**  heating average value calculated based on the day of the year and latitude of the measurement.  The uncertainty includes the potential error of ±20% in updraft velocity (\(w\)) and the standard error of the CCN concentration measurements.  The 25th, 50th, and 75th percentile in SWCF and CDNC obtained from the distribution of profiles generated by the \(ENT_{ML}\) analysis (section 2.5; Figures 1 and 2).
2.2. Aerosol Cloud Parcel Model (ACPM) Description

A 0-D ACPM was used in this study because (1) it enables the simulated results to be constrained by measured thermodynamic properties, aerosol number size distributions, and CCN spectra (Table 1; Figure S1 in the supporting information) and (2) the ACPM allows explicit comparison to in situ measurements of cloud droplet size distributions (rather than parameterized cloud droplet size distributions used in LES). The ACPM is based on a fixed-sectional approach to represent the dry particle size domain, with internally mixed chemical components. The model employs a dual-moment (number and mass) algorithm (Tzivion et al., 1987) to calculate particle growth from one size section to the next for nonevaporating compounds (viz., all components other than water) using an accommodation coefficient of 1.0 (Raatikainen et al., 2013). Liquid water is treated in a moving-section representation, similar to the approach of Jacobson et al. (1994), to account for evaporation and condensation of water in conditions of varying humidity. The model includes a dynamic scheme for activation of particles to cloud droplets (Seinfeld & Pandis, 2006) dependent on aerosol properties (Kohler, 1936; Petters & Kreidenweis, 2007). In subsaturated conditions (i.e., relatively humidity <100%), particles below the cloud base are considered to be in local equilibrium with water vapor pressure. A detailed description of the ACPM is presented in Russell and Seinfeld (1998) and Russell et al. (1999).

The criteria for selecting cases and the model inputs derived from variables, shown in Figure S1, are discussed in Text S1. The ACPM in-cloud lapse rate is adiabatic and ultimately compared to observations to assess entrainment (sections 2.4 and 2.5). SWCF is derived from the vertical profile of the in situ (CDP) and simulated (ACPM) cloud droplet size distribution and an asymmetric scattering parameter of 0.85 (see Text S2) (Bohren & Battan, 1980; Gerds et al., 2006; Hansen & Travis, 1974; Sanchez et al., 2017; Stephens, 1978). Coagulation, scavenging, and deposition of the aerosol were included in the ACPM, but their effects are negligible given the relatively short simulations used here (<1 hr) and their relatively low marine total aerosol particle concentrations (<1,000 cm^-3) and cloud LWC (<1 g m^-3). As stated previously, cloud droplet sizes were <20 μm diameter (for all but one case; RF10Thick; Table 1); therefore, autoconversion and accretion rates are not included in this simulation and exert a negligible effect on the simulated values of LWC and CDNC (Feingold et al., 2013).

2.3. ACPM Cloud Top Entrainment

The process of entrainment and its effect on a cloud can be described as the weighted downward mixing of above-cloud warm, dry air into the cloudy air below causing cloud droplets to evaporate. This air parcel subsequently mixes downward, diluting the cloud droplet concentration throughout the cloud, resulting in a reduction in CDNC compared to an adiabatic profile (Chen et al., 2011; Wang et al., 2003; Xue & Feingold, 2006). Five cases are available from this campaign to illustrate the effect of entrainment on SWCF in the SO (Table 1). These five cases are from four different RFs and are referred to by their flight number, with the two RF10 cases being differentiated by subscripts that describe the cloud thickness (RF10thin or RF10thick). Simulated cloud profiles are produced using the ACPM using adiabatic conditions, which are then compared to in situ observations and corrected for cloud top entrainment.

2.4. Entrainment Derived From LWC (ENTLWC)

The relative difference between simulated and observed values of LWC (or LWP) is related to the entrained fraction of above-cloud air. To account for the decrease in LWC associated with entrainment, the CDNC in the ACPM is decreased until the simulated (ACPM) and observed (CDP) LWC vertical profiles are the same. This approximation assumes inhomogeneous mixing, meaning only a fraction of the cloud droplets evaporate, whereas homogenous mixing implies all cloud droplets partially evaporate. Inhomogeneous mixing is consistent with past studies of well-mixed, nonprecipitating stratocumulus clouds (Brenguier et al., 2011; Burnet & Brenguier, 2007; Jia et al., 2019). Furthermore, homogeneous mixing would not sufficiently account for the observed decrease in SWCF and CDNC. In this study, evaporation of droplets due to inhomogeneous mixing is independent of droplet size; therefore, the volumetric mean diameter (Dv) does not change due to droplet evaporation.

2.5. Entrainment Derived From a Mixing Line (ENTML)

ENTML accounts for the subadiabatic LWC values due to cloud top entrainment with inhomogeneous mixing, derived from temperature, pressure, and qv measurements. ENTML, first used by Sanchez et al. (2017) and Calmer et al. (2019), generates a mixing line (a linear relationship between two conservative
variables; Figure 1) between \( q_v \) at the cloud base and above cloud, and measurement-derived equivalent potential temperature (\( \theta_E \)). \( q_v \) is equivalent to total water content (\( q_t \)), a conservative variable, in these subsaturated, nonprecipitating conditions. With the mixing line, the in-cloud \( q_t \) is a function of \( \theta_E \), which is derived from in-cloud temperature, pressure, and \( q_v \) (see Text S3 and Figure S3). The measured thermodynamic properties of the above-cloud air (shown as red circles in Figure 1) and cloud base air are then used to estimate subadiabatic LWC profile (Calmer et al., 2019; Sanchez et al., 2017). In this analysis, a line is made between a point at the cloud base and each point above the cloud that is identified as "entrained air" (red-bounded gray markers in Figures 1 and S4). Each line is used to calculate an entrainment-corrected vertical profile of CDNC, LWC, cloud droplet extinction, and integrated SWCF. The solid black line in Figures 1 and 2 is the 50th percentile (median) of the CDNC, LWC, and extinction, while the black dashed lines represent the 25th and 75th percentiles. See the Texts S3 and S4 for details on the measurements that represent the entrained air that are used to derive subadiabatic cloud LWC, CDNC, and \( D_v \).

3. Results

The observed LWC profiles in the SO are subadiabatic (Figures 2b and 2f) for all cases studied during the SOCRATES experiment. CDNCs are also always lower than the adiabatic simulations (Table 1; Figure 2) while \( D_v \) is consistent with adiabatic simulations, suggesting entrainment of dry air is reducing LWC via the evaporation of cloud droplets. These results are consistent with previous observations in the SO (Yum et al., 1998), the tropics (Rauber et al., 2007; Roberts et al., 2008), and midlatitudes (Calmer et al., 2019; Sanchez et al., 2017). Consequently, the SWCF and CDNCs above cloud base are significantly and consistently overestimated compared to adiabatic conditions (Table 1).

3.1. \textit{ENT}_{LWC}: Comparison of ACPM to Observations

To illustrate the impact of entrainment on cloud microphysical and optical properties, Figure 2 shows the vertical profiles of observed and simulated \( D_v \), LWC, CDNC, and cloud droplet extinction for RF12 and RF10\textsubscript{Thick} to illustrate two distinctly different cases for which entrainment corrections may be applied.
The adiabatic ACPM overestimates the observed LWC, CDNC, and cloud droplet extinction; however, $D_v$ remains remarkably similar to observed values, which supports evidence for inhomogeneous mixing (Brenguier et al., 2011; Burnet & Brenguier, 2007; Jia et al., 2019).

In-cloud level legs, which are straddled by vertical profiles for the RF13 case (Figure S6), reveal that LWC, $D_v$, and CDNC vary consistently together on scales over tens of kilometers. The tens of km-scale changes in the cloud base height (as well as cloud top height) are also shown by the HCR reflectivity during the RF13 in-cloud level leg (Figures S6 and S7). This large-scale horizontal variability in cloud properties and cloud base height is related to subcloud layer RH and temperature fields, which dictate lifting condensation levels, and is consistent with the findings of Wood and Taylor (2001). On length scales of a few km, updrafts are associated with higher CDNC and LWC compared to downdrafts, which is also consistent with the dilution of air parcels in downdrafts from cloud top entrainment (Wang et al., 2009; Yum et al., 2015). While $D_v$ follows a similar large-scale (>10 km; <10^{-2} s^{-1}) trend as LWC and CDNC, there is more than an order of magnitude less strength in variations of $D_v$ on smaller spatial scales (Figure S8). The similar strength in variations between LWC and CDNC at all length scales compared to $D_v$ further supports evidence of inhomogeneous mixing throughout the cloud layer—in that LWC is primarily controlled by the evaporation of cloud droplets rather than a shift in $D_v$. In spite of the variability in cloud base heights, simulated CDNC and extinction reproduce observed values, when accounting for entrainment with $\text{ENT}_{\text{LWC}}$ (section 2.4).

**Figure 2.** Observed and simulated cloud properties for the RF12 (a–d) and RF10\text{Thick} (e–h) case. The observed properties were obtained from the CDP measurements, except for the extinction, which was calculated from the observed CDP distributions. The $\text{ENT}_{\text{LWC}}$ and $\text{ENT}_{\text{ML}}$, $D_v$ profile is identical to the adiabatic profile due to the assumption of inhomogeneous entrainment. By definition, the LWC for $\text{ENT}_{\text{LWC}}$ is the same as what was observed, so there is no $\text{ENT}_{\text{LWC}}$ line on the LWC profile. The $\text{ENT}_{\text{ML}}$ profile is dependent on the properties of the entrained air (red circles in Figure 2). The solid black line is the median LWC and CDNC from the distribution of profiles derived from the $\text{ENT}_{\text{ML}}$ analysis (section 2.5), and the dashed lines represent the 25th and 75th percentiles and correspond to the mixing lines shown in Figure 1. Vertical profiles for RF02, RF10\text{Thin}, and RF13 are shown in Figure S5.
Table 1 shows the difference in SWCFs comparing observations to the adiabatic model (5SWCF\textsubscript{adiabatic}) and accounting for entrainment using LWC (5SWCF\textsubscript{ENT,LWC}) and mixing line (5SWCF\textsubscript{ENT,ML}) corrections. A difference of 0 W m\(^{-2}\) implies consistency between observations and the entrainment method. 5SWCF\textsubscript{adiabatic} represents the upper limit and ranged from 38.4 ± 8.9 to 93.4 ± 17.1 W m\(^{-2}\), while 5SWCF\textsubscript{ENT,LWC} are significantly lower, between 5.3 ± 16.2 and 32.1 ± 14.2 W m\(^{-2}\). As LWC in the observed and ENT\textsubscript{LWC} vertical profiles are, by definition, the same, 5SWCF\textsubscript{ENT,LWC} is greater than zero because of narrower droplet spectrum width and greater CDNC in the simulations. Compared to the adiabatic simulations, CDP-LWP was reduced by 28–58% as a result of cloud top entrainment. For these five cases, the range in SWCF associated with 28% and 58% reduction in LWP compared to an adiabatic LWP amounts to 53 W m\(^{-2}\) (Figure S8).

This implies that variability in entrainment exerts an uncertainty in radiative forcing of individual clouds up to 50 W m\(^{-2}\).

The ACPM was used to further quantify and compare the sensitivity of SWCF and CDNC to changes in aerosol (N\textsubscript{p}), updraft (q\textsubscript{t}), and LWP related to entrainment (Table 1). In these sensitivity calculations, N\textsubscript{p} and q\textsubscript{t} distributions were evaluated at half and double their observed values (i.e., 0.5N\textsubscript{p} and 2N\textsubscript{p}; 0.5q\textsubscript{t} and 2q\textsubscript{t}). Table 1 shows that SWCF is two to four times more sensitivity to N\textsubscript{p} than q\textsubscript{t} over the SO. Similarly, CDNC was approximately 1–3 times more sensitive to N\textsubscript{p} than q\textsubscript{t}. This indicates that clouds in the SO form in an aerosol-limited regime, which is characteristic of the MBL (Reutter et al., 2009; Sanchez et al., 2016). Table 1 shows that the SWCF sensitivity to entrainment and subsequent reduction of LWP are 2–20 times those related to N\textsubscript{p} and q\textsubscript{t}. CDNC is approximately twice as sensitive to entrainment compared to N\textsubscript{p} and up to four times more sensitive than the updraft. This highlights the first-order role of entrainment in determining cloud microphysical properties.

The SWCF for RF02 and RF10\textsubscript{Thick} illustrate limiting cases in this study as they are the most and least sensitive, respectively, to both N\textsubscript{p} and entrainment. The RF02 SWCF case shows the greatest sensitivity to N\textsubscript{p}, and entrainment because it has the lowest CDNC concentration and it is the thinnest cloud of the five cases (5SWCF\textsubscript{ENT,LWC} = 32.1 ± 14.2 W m\(^{-2}\); Table 1). In contrast, the impact of entrainment on RF10\textsubscript{Thick} is lower than the other cases in this study because the cloud is significantly thicker (Table 1), and the cloud SWCF is less sensitive to changes in optical thickness when the optical thickness is high (Sanchez et al., 2016) (Figure S8). The RF10\textsubscript{Thick} case is also unique relative to the other cases because it has a large Aitken mode, as well as the highest CCN concentration (Figure S1c and Table 1). The RF10\textsubscript{Thick} CDNC is more sensitive to N\textsubscript{p} than q\textsubscript{t} (aerosol-limited case; Reutter et al., 2009); however, it is relatively more sensitive to q\textsubscript{t} compared to the other cases in this study because of the high Aitken mode particle concentration (similar to findings by Sanchez et al., 2016).
4. Discussion and Implications

This study uses in situ measurements from the SOCRATES, collected on the NSF/NCAR HIAPER Gulfstream V (GV) during the austral summer. Measurements were used to assess the impact of cloud top entrainment by combining in situ observations of CCN spectra, aerosol number size distributions, and updrafts at cloud base to initialize an aerosol-cloud parcel model (ACPM) and compare simulated cloud microphysical and optical properties to those derived by a CDP. Differences in SWCF between adiabatic simulations and subadiabatic observations of a nonprecipitating stratocumulus cloud layer in the SO resulted in values as high as ~100 W m$^{-2}$. SWCF was simulated to within 16 W m$^{-2}$ using CDP-derived entrainment ($ENT_{LWC}$) for four of the five cases. In one of the cases (RF02), the difference between $ENT_{LWC}$ and observed SWCF is relatively large due to the higher sensitivity of SWCF to particle concentrations. This case was also the thinnest cloud in this study, and SWCF sensitivity to entrainment is higher for such conditions (Figure S8).

Even when accounting for the impacts of entrainment on cloud optical properties, the entrainment-corrected LWP varied between 28% and 58%, which translates to a range of 53 W m$^{-2}$ in SWCF related to variability in entrainment for the cases studied here. For comparison, the sensitivity of SWCF due to entrainment is between 2 and 20 times the sensitivities of SWCF related to particle concentration and updraft—thus emphasizing the importance of accurately accounting for entrainment in climate models.

$ENT_{ML}$ (mixing-line approach) was simulated to within 15 W m$^{-2}$ except for two cases. $ENT_{ML}$ utilizes the linear relationship of two conservative variables (total water content [$qt$] and $θ_E$) to derive in-cloud $qt$ from the measured in-cloud $θ_E$ to calculate the entrainment-induced reduction in LWP. This method was first shown in pristine and polluted midlatitude clouds (Calmer et al., 2019; Sanchez et al., 2017) and is now reproduced here with in situ CDP measurements for clouds in the SO. $ENT_{ML}$ only requires measurements of pressure, temperature, and relative humidity (PTU) in and above the cloud as well as the vertical extent of the cloud to estimate the subadiabatic LWC profile. $ENT_{LWC}$ is certainly a more accurate method to account for entrainment but also requires the deployment of a CDP (or similar instrument to measure CDNC and LWC), which is not nearly as ubiquitous as PTU measurements. PTU measurements can be obtained with numerous airborne platforms (i.e., radiosondes, aircraft, and UAVs) as long as the vertical resolution is high enough to capture the temperature and water vapor immediately above the cloud layer (within ~10 to 100 m).

$ENT_{ML}$ approaches an adiabatic profile (i.e., no entrainment correction), if $θ_E$ remains constant in the cloud while $qt$ actually varies. Nonetheless, $ENT_{ML}$ could be used to improve GCM parameterizations, such as multivariate PDFs representing subgrid-scale moisture, temperature, and vertical velocity to derive cloud microphysical parameters across multiple cloud regimes (Guo et al., 2010). Reducing the uncertainty of cloud optical properties in GCM models remains a fundamental challenge for correctly representing the Earth’s energy budget (Bretherton et al., 2013; Briant et al., 2019; de Szoeke et al., 2010; Hourdin et al., 2015). For decades now, the IPCC has reported an uncertainty in the anthropogenic-induced cloud forcing of approximately 1.5 W m$^{-2}$, which represents nearly 75% of the uncertainty related to the total anthropogenic radiative forcing (IPCC, 2014). This study in the SO extends beyond earlier studies in the midlatitudes (Calmer et al., 2019; Sanchez et al., 2016) to show that the reduction in LWP owing to entrainment is a major contributor to uncertainty in SWCF for stratocumulus clouds globally. Improving the parameterizations of entrainment using the mixing line techniques presented here has the potential to significantly improve GCM representation of cloud optical properties and, ultimately, the Earth’s energy budget.

Conflict of Interest

The authors have no conflicts of interest to disclose.

Data Availability Statement

SOCRATES data are available through the following EOL UCAR repository: https://data.eol.ucar.edu/mas-ter_lists/generated/socrates/ (Sanchez & Roberts, 2018; UCAR/NCAR-Earth Observing Laboratory, 2018; UCAR/NCAR-Earth Observing Laboratory, 2019a; UCAR/NCAR-Earth Observing Laboratory, 2019b). The VCSEL water vapor QC/QA data can be accessed online (at https://data.eol.ucar.edu/dataset/552.052).
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