INVESTIGATING CXOU J163802.6−471358: A NEW PULSAR WIND NEBULA IN THE NORMA REGION?

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Received 2014 January 20; accepted 2014 April 18; published 2014 May 12

ABSTRACT

We present the first analysis of the extended source CXOU J163802.6−471358, which was discovered serendipitously during the Chandra X-ray survey of the Norma region of the Galactic spiral arms. The X-ray source exhibits a cometary appearance with a point source and an extended tail region. The complete source spectrum is fitted well with an absorbed power law model and jointly fitting the Chandra spectrum of the full source with one obtained from an archived XMM-Newton observation results in best fit parameters $N_{\text{H}} = 1.5^{+0.7}_{-0.5} \times 10^{23}$ cm$^{-2}$ and $\Gamma = 1.1^{+0.7}_{-0.6}$ (90% confidence uncertainties). The unabsorbed luminosity of the full source is then $L_X \sim 4.8 \times 10^{33} d_{10}^{-2}$ erg s$^{-1}$ with a distance of 10 kpc a lower bound inferred from the large column density. The radio counterpart found for the source using data from the Molonglo Galactic Plane Survey epoch-2 shows an elongated tail offset from the X-ray emission. No infrared counterpart was found. The results are consistent with the source being a previously unknown pulsar driving a bow shock through the ambient medium.

Key words: acceleration of particles – pulsars: general – radio continuum: general – stars: individual (CXOU J163802.6−471358) – X-rays: general

Online-only material: color figures

1. INTRODUCTION

Most pulsars are discovered in radio surveys by their pulsation properties. When they are later resolved in X-rays with high resolution telescopes such as the Chandra X-Ray Observatory, complex structures associated with the pulsar are discovered. Since the pulsar loses rotational energy through a wind of ultra-relativistic particles, the non-thermal synchrotron emission that arises will be observable across the electromagnetic spectrum in the form of a pulsar wind nebula (PWN; see Gaensler & Slane 2006 and Kirk et al. 2009 for a review). All pulsars are expected to be associated with a PWN (Kargaltsev & Pavlov 2008), but only a fraction of them are detected at current sensitivity. The PWNe that have been identified exhibit a morphology dependent on the ambient medium, and on pulsar properties such as age, speed, and magnetic field strength. The high sensitivity and arcsecond angular resolution of the Chandra telescope makes it possible to study the X-ray morphology of these sometimes weak extended sources as well as their spectral properties. For a pulsar moving supersonically through the ambient medium, the PWN will be confined by ram pressure and display a cometary structure in X-rays. This type is generally referred to as a bowshock-tail PWN, where “The Mouse” PWN (Gaensler et al. 2004) is a typical example.

An X-ray survey of a $2^\circ \times 0.8$ region of the Norma spiral arm was done during summer 2011. Twenty-seven observations, each of $\sim 20$ ks duration, were performed with the Chandra X-Ray Observatory for an in depth study of the region. The primary goal of the survey was to find new high-mass X-ray binaries (HMXBs) in order to study their evolution. The Norma region was chosen because the hard X-ray imaging of the Galactic plane done by the INTEGRAL Gamma-Ray Astrophysics Laboratory revealed that the Norma region contains the largest number of HMXBs and OB associations after the Galactic center (Fornasini et al. 2014). Around 1500 sources were found in the Chandra survey and a catalog of these sources is presented in another paper (Fornasini et al. 2014). Out of these 1500 sources at least 6 are extended. Four of these have been previously detected and categorized as supernova remnants (SNRs) and are mentioned in SNR catalogs such as the Green catalog (Green 2009) and the new high energy Galactic SNR catalog (Ferrand & Safi-Harb 2012). One of the other two new sources is treated in this paper.

In Section 2 we present the data that was utilized in this investigation. In Sections 3.1 and 3.2 we present the methods used and the results that were achieved based on the X-ray data available. Section 3.3 describes the attempt to identify counterparts to the X-ray emission using infrared and radio data. Finally, in Section 4 we evaluate the results in order to identify the source type.

2. OBSERVATIONS

The extended source CXOU J163802.6−471358 was serendipitously covered by two Chandra exposures with ObsIDs 12519 and 12520 (19.3 ks and 19.0 ks exposures, respectively). Both were performed on 2011 June 13 using the ACIS-I instrument in VFAINT mode, and both had the source positioned off-axis, more so in ObsID 12520, where it is located at $\theta \sim 8.4$ compared to $\theta \sim 3.8$ for ObsID 12519.

All Chandra data preparation was performed using the Chandra Interactive Analysis of Observations (CIAO) package version 4.4. The event files to be used for imaging analysis were filtered on energy so as to only include energies in the range of 0.5–8.0 keV. In order to increase the signal to noise ratio, the two observations were merged using the CIAO routine merge Obs, providing a merged exposure corrected image of the source. The spectral analysis, on the other hand, was done using the unfiltered level 2 event files.

Contaminating point sources of significance larger than 4$\sigma$ as listed in the Norma region point source catalog (Fornasini et al. 2014) were removed prior to the analysis.
CXOU J163802.6−471358 was found to be covered by an archived XMM-Newton observation from 2005 August 19–20, and the source has been cataloged as 3XMM J163802.6−471357. In this work, we include the data from this observation (ObsID 0307170201), which consists of a 99.5 ks pointing during which CXOU J163802.6−471358 was 11°5 off-axis. Including the XMM-Newton data allows us to compare to previous measurements and to improve the constraints on the spectral parameters.

The source field has also been covered by several surveys over multiple wavelengths, as both radio data from the Molonglo Galactic Plane Survey epoch 2 (MGPS-2; Murphy et al. 2007) and infrared observations from the Spitzer Space Telescope5 and the Visible and Infrared Survey Telescope for Astronomy (VISTA)6 are available. We make use of these survey images in order to look for counterparts to the X-ray source.

3. ANALYSIS AND RESULTS

A quick examination of the merged Chandra image (Figure 1) reveals a complex structure of CXOU J163802.6−471358 featuring two distinct components: a point-like part with detection significance 38.9σ and a diffuse part with significance 9.9σ, both calculated for the energy range 2.0–8.0 keV. This energy range was chosen since no source counts were detected below 2.0 keV, presumably as a result of the high column density toward the Norma region. The point-like source is positioned at (α, δ)2000 = (16°38′02″.7, −47°13′58″/4) in equatorial coordinates with a position error of 0.6 (at 95% confidence; Fornasini et al. 2014) and the source significance was calculated based on the events falling inside a circular region centered on that position, having a radius of 3.5′ (the white circular region in Figure 1). The significance of the diffuse emission was calculated from the events falling inside the rectangular region shown in Figure 1. Since the significance is calculated based on the summed counts in both exposures, the emission from the point-like source is somewhat blended with the diffuse extended emission, and we chose the rectangular region so that it included as many counts as possible without including too many of the point source photons. The diffuse part is tail-like (referred to as “the tail” hereafter) and extends from the point-like source ∼40″ to the north.

Furthermore, when smoothing the image, two diffuse features emerge. One of them is only visible in ObsID 12520 (marked with a dashed ellipse in Figure 1), and its significance is right below the limit for positive detection at 2.9σ, and stretches ∼19.5″ to the West. The other one is possibly observed in ObsID 12519 below the significance threshold at 2.4σ, but is not visible in the merged image (Figure 1), since it is blended with the point-spread function (PSF) of observation 12520.

3.1. Imaging Analysis

An exposure corrected, full resolution (0.492″ × 0.492″), merged Chandra image of CXOU J163802.6−471358 for the energy range 2.0–8.0 keV is shown in the top left panel of Figure 2. The source is hard, peaking at an energy of ∼5 keV, and in order to examine the morphology of the source with energy, a merged, exposure corrected image was additionally created for three narrow energy bands, representing soft, medium, and hard X-ray energies. The energy bands were chosen so that a similar number of source net counts (∼100) were present in each band. Since no source counts were detected below 2.0 keV, the selected energy ranges were 2.0–4.0 keV (soft), 4.0–5.5 keV (medium), and 5.5–8.0 keV (hard), as shown in Figure 2.

The spatial extension of the source was examined further through a surface brightness profile constructed using aperture slices taken along the symmetry axis. Only observation 12519 was used in this step, since the tail and point-like source were expected to be somewhat blended in observation 12520, caused by the location of the source at a large off-axis angle. The resulting profile can be seen in Figure 3 with the inset image showing the extraction regions used. Each rectangular region measures 30″ × 4′ across the tail section and to the south of the point-like source, whereas the regions lying on top of the point-like source measure 30″ × 2′.

In order to determine whether the main component is in fact a point source, a Chandra PSF was simulated at the source position, and superposed on the profile. The PSF was generated using the Chandra Ray-Viewing (ChaRT) software7 and simulated using the MARX simulator version 5.0.

3.2. Spectral Analysis

In order to look for spectral variations across the Chandra source, we extracted spectra from the two source components (the point source and the tail) along with a spectrum of the full source area. The extraction regions were defined as shown in Figure 1, and all spectra and the associated RMF and ARF files were created using the CIAO tool specextract. Events from both observations 12519 and 12520 were extracted for each region, and combined to remediate the sparse number of counts in each exposure. One background spectrum was also extracted for each exposure from a source-vacant area on the chip, and using the Chandra PSF generated using the ChaRT software7 and simulated using the MARX simulator version 5.0.

5 http://irsa.ipac.caltech.edu/data/SPITZER/GLIMPSE/
6 http://horus.roe.ac.uk/vsa/index.html
7 http://cxc.harvard.edu/chart/runchart.html
the $N_H$ parameter for the three tied together and the photon index varying freely. This was done assuming that the column density $N_H$ does not vary across the source, which is in good agreement with the fact that the best fit value of the $N_H$ parameter from individual fits is the same within 90% confidence. The resulting best fit parameters for the three regions can be seen in the top half of Table 1.

After fitting the grouped spectra, we examined whether using Cash statistics (Cash 1979) on the ungrouped spectra gave a significantly different result. The fit parameters were consistent within 90% confidence and the results based on the grouped spectra were therefore used in the further analysis.

Given the generally poor photon statistics, a hardness ratio was additionally calculated for each source component using the general formula $HR = (H - S)/(H + S)$, where $H$ is the net count number in the hard energy band 4.5–8.0 keV and $S$ is the net count number in the soft energy band 2.0–4.5 keV. The energy ranges were again chosen so that approximately the same number of source counts from the full source field reside in each band. The calculation resulted in a hardness ratio with 1σ errors of 0.16 ± 0.10 and 0.08 ± 0.21 for the point source and the tail component, respectively, meaning that there is no significant difference between the hardness of the two.

We made an XMM-Newton spectrum for the PWN candidate using the Scientific Analysis Software (SAS) version 13.0.1. We produced new event files using ep_prob and emproc and filtered them according to the recommended procedures. Based on the 10–12 keV PN light curve, we removed events that occurred during background flares. Such flares were prevalent during this observation, and we were only able to use approximately 8 See http://xmm.esac.esa.int/sas/current/documentation/threads/.

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**Figure 2.** PWN candidate is shown here in four different energy ranges. All four images are exposure corrected and unbinned with pixel size $\sim 0\arcsec.492 \times 0\arcsec.492$. Panels (b)-(d) have additionally been smoothed using a Gaussian kernel of size $\sigma = 2\arcsec.5$. North is to the right and east is up.

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**Table 1**

| Region                | $N_H$ ($10^{20}$ cm$^{-2}$) | $\Gamma$    | Goodness of Fit ($\chi^2$/d.o.f.) | $F_{12-10keV}^{\text{obs}}$ (10$^{-13}$ erg cm$^{-2}$ s$^{-1}$) | Counts No. |
|-----------------------|-----------------------------|-------------|-----------------------------------|-------------------------------------------------|-------------|
| 1. Full source        | $13^{+6}_{-5}$              | $0.7^{+1.0}_{-0.9}$ | 60.5/54                           | $3.2^{+0.9}_{-1.0}$                               | 356         |
| 2. Point source       | ...                         | $0.5^{+1.0}_{-0.9}$ | ...                              | $2.5^{+0.7}_{-0.8}$                               | 180         |
| 3. Tail               | ...                         | $1.3^{+1.4}_{-1.2}$ | ...                              | $0.8^{+0.4}_{-0.5}$                               | 110         |
| XMM full source       | $15^{+11}_{-5}$             | $1.4^{+0.9}_{-0.8}$ | 38.4/37                           | $2.6^{+0.3}_{-0.8}$                               | 306$^b$     |
| XMM+Chandra           | $15^{+11}_{-5}$             | $1.1^{+0.7}_{-0.6}$ | 77.8/70                           | $2.6^{+0.4}_{-0.5}$                               | ...         |

**Notes.** Uncertainties are at the 90% confidence level.

$^a$ Based on the abundance values given in Wilms et al. (2000).

$^b$ Number of counts quoted is for the PN spectrum.

$^c$ The absorbed flux value quoted here is calculated based on the PN spectrum. The value calculated from the ACIS spectrum and the MOS2 spectrum are consistent with this value, whereas there is a slight discrepancy between them and the MOS1 value.
half of the time, leaving exposure times of 48.5 ks for PN and 54.2 ks for each MOS unit. The angular resolution of XMM-Newton does not allow us to separate the point source from the extended emission, and, for all three instruments (PN, MOS1, and MOS2), we extracted a spectrum from a 30′′ radius circle, which includes both the point source and the tail. In order to subtract the background, we used nearby rectangular regions that do not include sources. We re-binned the spectra to require a signal-to-noise ratio of at least 2.7 in each bin (except for the highest energy bin). We obtained source count rates of 6.3 × 10⁻³, 3.2 × 10⁻³, and 2.2 × 10⁻³ counts s⁻¹ for PN, MOS1, and MOS2, respectively.

The three spectra from the three instruments (PN, MOS1, and MOS2) were subsequently loaded into XSPEC and fitted simultaneously using the same absorbed power-law model as used for the Chandra ACIS spectra. The column density N_H was tied together and all parameters including an inter-instrumental normalization constant were allowed to vary freely. The best fit parameters resulting from the fit can be seen in the bottom part of Table 1.

The Chandra and XMM-Newton best fit parameters were found consistent within 90% confidence, and based on this, the Chandra full source spectrum and the three XMM-Newton spectra were fitted simultaneously following the same fitting procedure as previously. The fitted full source spectra and their residuals can be seen in Figure 4, and the best fit parameters are listed in Table 1.

### 3.3. Radio and Infrared Imaging

A large part of the Southern Galactic Plane is covered by the MGPS-2, which includes the Norma region and source CXOU J163802.6−471358. This provides us with the means to search for a radio counterpart to the X-ray source.

The MGPS-2 was done using the Molonglo Observatory Synthesis Telescope at a frequency of 843 MHz and a restoring beam size of 45 × 45 csc|δ| arcsec², where δ is the declination.

The particular MGPS-2 mosaic image covering the source region was examined, and a radio structure was found coinciding with the X-ray peak position within the resolution of the observation. An image of the X-ray emission with an overlay of radio contours can be seen in Figure 5, with the beam shape and size indicated. The radio structure seems to be an unresolved radio peak toward the north. An estimate of the peak brightness of the radio structure is found as the value of the brightest pixel at 65.5 mJy beam⁻¹.

We furthermore found through archival searches that the source area has been covered by both the Spitzer Space Telescope and the VISTA, providing images in the mid- and near-infrared. The Spitzer survey, called the Galactic Legacy Infrared Mid-Plane Survey Extraordinaire (GLIMPSE; Benjamin et al. 2003), covers a large part of the Galactic disk, including the source CXOU J163802.6−471358. The Spitzer IRAC instrument collected images in four bands (3.6, 4.5, 5.8, and 8.0 μm) and a 4.5 μm image of the source region with an overlay of Chandra X-ray contours is shown as the middle panel of Figure 6.

The VISTA images were obtained from the VISTA Science Archive using the VISTA Variables in the Via Lactea database (Saito et al. 2012). The four filters ZYJHK_s, corresponding to 0.88, 1.02, 1.25, 1.65, 2.15 μm respectively, were all examined and the K_s band is shown as the bottom panel of Figure 6. The limiting magnitude for the K_s band is ~18 mag (Minniti et al. 2010).

### 4. DISCUSSION

The source CXOU J163802.6−471358 is visibly extended across all the energy ranges presented in Figure 2. A comparison between the radial profile of the source and a simulated Chandra PSF (Figure 3) confirms that the source consists of two components, a point source and an extended tail. Both components are detected to a high significance of 38.9σ and 9.9σ, respectively.

Based on the extension of the source, a few possibilities were considered for the source type, an active galactic nucleus (AGN) with a jet, a low-mass X-ray binary (LMXB) with a jet, and a bow-shock tail PWN. The results presented in the previous section make us prefer the latter for the following reasons.
Figure 5. Shown here is the exposure corrected merged image of CXOU J163802.6−471358 in grayscale, smoothed with a Gaussian kernel of size $\sigma = 2''$. On top is an overlay of radio contours from the MGPS-2 showing 7–70 mJy beam$^{-1}$ in 10 square root steps. The beam shape and size of $45 \times 45 \csc[\delta]$ arcsec$^2$ is shown in the top right corner of the image.

Figure 6. Infrared view of the source field. Top plot: Chandra image with Chandra contours. The panel measures $\approx 1'0 \times 0'4$. Middle plot: infrared 4.5 $\mu$m obtained from the GLIMPSE survey. Bottom panel: infrared 2.15 $\mu$m from VISTA. The circle indicates the position of the Chandra point source and its 0'6 position error, which includes the absolute astrometric accuracy. The nearest infrared source is located 1.7'' away in the VISTA survey. The square cut-out in the bottom panel measures 5'' × 5''.
The full source spectrum (point source + tail) is fitted well with an absorbed power-law, pointing toward a synchrotron origin of the emission. The column density at \( N_H = 1.5 \times 10^{23} \text{cm}^{-2} \), resulting from the joint \textit{Chandra} and XMM-Newton fit, is rather high and indicates that the source is located in the far Norma arm. Comparing to another source in the same area, PWN HESS J1640-465 (Lemiere et al. 2009), it is apparent that the Norma arm is a region of high absorption since this source exhibits an X-ray spectrum of comparable column density (\( N_H = 1.4 \times 10^{23} \text{cm}^{-2} \)). This is to be expected, since the Norma region is known to host a large number of molecular clouds (Dame et al. 2001).

The distance extinction in the line of sight toward the Norma region was found through IRSA\(^9\) to be \( A_V = 64 \text{ mag} \) (Schlafly & Finkbeiner 2011). Using this together with the relation \( N_H/A_V = 2.5^{+0.4}_{-0.3} \times 10^{21} \text{ cm}^{-2} \text{ mag}^{-1} \) (Watson 2011) results in an expected column density of \( 1.4 \times 10^{23} \text{ cm}^{-2} \) for the particular region of this source. This makes the column density obtained from fitting a reasonable value for a source in the far Norma arm, but it cannot be excluded that the source could be extragalactic.

The distance to the HESS source was estimated based on measurements of H\(\alpha\) absorption features toward the system, by comparing the derived velocity to the Galactic rotation curve (Lemiere et al. 2009). The comparable column density of the two sources, implies that the distance estimated for the HESS source (~10 kpc) can be accepted as a lower bound on the distance to CXOU J163802.6–471358 and is used as such for the remainder of this discussion.

The photon index of \( \Gamma = 1.1^{+0.7}_{-0.6} \) obtained from the joint \textit{Chandra} and XMM-Newton fit is consistent with what is expected for a pulsar+PWN spectrum, where \( \Gamma_{\text{PSR}} \) and \( \Gamma_{\text{PWN}}\) usually lie in the range \( 1 \lesssim \Gamma \lesssim 2 \) (Kargaltsev & Pavlov 2008). This value implies a somewhat harder spectrum than what is generally seen from an AGN (Mainieri et al. 2007) and the ACIS spectrum obtained for the point source is even harder (\( \Gamma = 0.5^{+1.0}_{-0.9} \)).

Comparing the flux obtained from the \textit{Chandra} observations, and the XMM-Newton observation done 5 yr previously \( (F_{\gamma, \text{Chandra}} = 3.2^{+0.9}_{-1.0} \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1} \) and \( F_{\gamma, \text{XMM}} = 2.6^{+0.3}_{-0.8} \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1} \), respectively), no long timescale variability is seen. Since an AGN is expected to exhibit the strongest long timescale variability in the soft X-ray regime (<12 keV) (Chitnis et al. 2009), this result decreases the likeliness of the source being an AGN, unless we just happened to catch the source when it was at comparable flux levels.

An LMXB is also expected to be strongly variable which makes this hypothesis less likely as well. The absence of X-ray variability on long timescales supports the PWN hypothesis, where usually only short timescale variations at the pulse period are expected.

We attempted a pulsation search with XMM-Newton, making use of the timing resolution of the PN instrument (~73 ms); unfortunately the poor photon statistics combined with the high background during the observation did not allow us to place any constraints on a possible pulsar period.

Since the flux values from the two observation epochs (\textit{Chandra} and XMM-Newton) are consistent, we used the unabsoled flux obtained from the joint fit together with the distance given above in order to calculate a luminosity of the full source (point source + tail). This gives a luminosity in the 2–10 keV energy range of \( \sim 4.8 \times 10^{33}d_{150}' \text{ erg s}^{-1} \) which is in good agreement with the values measured for several pulsars (Tables 1 and 2 in Kargaltsev & Pavlov 2008).

A radio counterpart to the X-ray source CXOU J163802.6–471358 was found from MGPS-2 data, and revealed a radio trail aligned with the X-ray trail (Figure 5), which indicates that there is a connection between the two. The synchrotron origin of the emission suggested by the spectrum of the source is further supported by the fact that the radio emission stretches further to the north than does the X-ray emission. This could be the result of synchrotron aging effects, meaning that the low energy emitting electrons live longer than the high energy emitting electrons and therefore are able to travel further away from the origin of the emission (the point source). If this is indeed a radio counterpart to the X-ray source, the likelihood of the source being an LMXB is small, since only a very rare type of microquasar, such as 1E 1740.7–2942 or GRS 1758–258 (Mirabel et al. 1993), with persistent radio jets would be likely to be found with extended radio emission. However, even in this LMXB scenario, the properties of CXOU J163802.6–471358 are not a good match since the persistent radio jets are double-sided and the X-ray emission from most LMXBs is highly variable on a wide range of timescales, making it extremely unlikely that the \textit{Chandra} and XMM-Newton observations, separated by 6 yr, would have the same or nearly the same flux. Some exceptions do exist where the LMXBs are in a quiescent phase, only accreting at a low level and therefore exhibiting long periods of relatively stable luminosity, which typically lies in the range \( L_X = 10^{30}–10^{33} \text{ erg s}^{-1} \) (Garcia et al. 2001). The luminosity of \( L_X \sim 4.8 \times 10^{33}d_{150}' \text{ erg s}^{-1} \) for CXOU J163802.6–471358 indicates that if this was an LMXB in its quiescent phase, it would be at the high end of the expected luminosity range for such sources, where significant variability is known to occur (see Cackett et al. 2011 and Bradley et al. 2007 for examples).

An offset between the radio and X-ray peak of ~40″ is observed, and given the absolute astrometric offset of <2″ for the radio survey (Murphy et al. 2007) and the 0′.6 positional uncertainty for the \textit{Chandra} source, the offset is real, assuming that the radio emission is in fact a counterpart to the X-ray source. Such an offset is seen for some PWNe driven by pulsars with high kick velocity (PSR J1509−5850 in Kargaltsev et al. 2008, IGR J11014−6103 in Pavan et al. 2014). Such an offset is not expected for an AGN where the nucleus should be a static emitting source with the radio peak coincident with the X-ray peak.

The nearest IR source is located 1′′7 from the X-ray source (Figure 6), and taking the positional error of 0′.6 (white circle in Figure 6) on the \textit{Chandra} source position into account, the IR source can be excluded as a counterpart to the X-ray source. This is consistent with the PWN interpretation, where the absence of an IR counterpart is expected. However, it is not an extremely strong constraint since any counterpart would be hidden by the large extinction toward the Norma region.

We exclude the possibility of the source being an LMXB based on the above, and we argue that an AGN nature of the source is unlikely. On the other hand, no results of this analysis challenge the hypothesis of the point source being a high velocity pulsar driving a ram pressured PWN, which makes us believe that this is indeed a previously unknown bowshock-tail PWN.

Based on the bowshock-tail interpretation of the source, we can use the empirical relation log \( L_{\gamma,(2–10keV)} = 1.34 \log E - 15.34 \) given in Possenti et al. (2002) to calculate the spin-down luminosity of the pulsar from the total

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\(^{9}\) http://irsa.ipac.caltech.edu/applications/DUST/
unabsorbed source luminosity given above, which results in an estimate of the spin-down luminosity of $E = 4 \times 10^{36} \text{ erg s}^{-1}$. The estimate is highly uncertain due to the large scatter in the above relation, and could be off by a factor of 10. However, the value seems reasonable for a bowshock tail PWN (see Table 1 in Kargaltsev & Pavlov 2008), and is also close to the value found through radio timing observations of the pulsar PSR B1757−24 ($E \sim 2.6 \times 10^{36} \text{ erg s}^{-1}$; Kaspi et al. 2001), which exhibits a very similar X-ray structure to what is observed for this source. We do not present other estimated values for the pulsar spin parameters, since a careful examination of the relevant calculations reveals that the results vary a great deal depending on the values assumed.

Using the 10 kpc value as a lower bound on the distance to the source, and using the angular length of the X-ray tail of $\sim 40''$, the projected length of the tail is found to be $l_{\perp} \sim 1.9d_{10} \text{ pc}$. The true length of the X-ray tail could be significantly longer, since we did not detect any significant change in tail length with energy, or significant softening of the spectrum with increasing distance from the pulsar candidate, which would otherwise be expected for a synchrotron emitting source. This might be due to the poor photon statistics that would lead to detection of only the brightest part of the tail where no effects of synchrotron aging would be detectable.

Accepting the pulsar interpretation, makes it even more interesting to determine whether the jet-like feature perpendicular to the trail of the source (marked by the dashed ellipse in Figure 1) is real or not. Associations could be drawn toward other pulsars showing evidence of a jet protruding almost perpendicular to the velocity vector of the pulsar (i.e., IGR J11014−6103 in Pavan et al. 2014 and B2224+65 in Johnson & Wang 2010).

In order to prove or disprove the claim that the point source is a pulsar and that the extended emission is a bowshock tail created by ram pressure, we need longer exposure with Chandra and better resolution radio data. This would also provide us with the means of resolving the diffuse jet-like feature. A search for pulsations from the point source in radio or in X-ray using XMM-Newton, or Chandra in timing mode, would be very beneficial, and would provide the means to definitively determine the nature of the source. Making use of the increased resolution at hard X-ray energies provided by the Nuclear Spectroscopic Telescope Array we could possibly provide better constraints on the spectrum of this hard source.

The Dark Cosmology Centre is funded by the DNRF. J.A.T. acknowledges partial support from NASA through Chandra Award GO1-12068A issued by the Chandra X-Ray Observatory Center, which is operated by the Smithsonian Astrophysical Observatory for and on behalf of NASA under contract NAS8-03060. This work is based in part on observations made with the Spitzer Space Telescope, which is operated by the Jet Propulsion Laboratory, California Institute of Technology under a contract with NASA.

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