Optical Measurements of an Atmospheric Pressure Microplasma Jet Aiming Surface Treatment

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Abstract In this work stable non-thermal ac high voltage atmospheric pressure microplasma jet (APMJ) device was used for optical and electrical characterizations. It enables the generation of low power (~5W) microplasma jet at frequency of 60Hz. The jet has a visible radial diameter of approximately 1.5 mm. Optical emission spectroscopy was used as a diagnostic tool to determine the gas discharge parameters as the modes temperature. The rotational temperature of OH radicals at the exit nozzle varies from 325 to 525K for different gases where the electrical input power ranged from 3 to 10W. Both the electronic (0.5 – 0.7eV) and vibrational temperatures (0.35 – 0.58eV) were estimated at the same power conditions for Ar, He, N₂, air and O₂ + 1%Ar flow rates. The highly reactive species as OH, O, N₂⁺ and the energetic photons produced between the electrodes extent along the plasma plume, both in radial and axial direction from the exit of the APMJ.

Keywords Microplasma jet, Gas temperature, Optical emission spectroscopy

1. Introduction

The generation of atmospheric pressure microplasma jets (APMJ) using an alternating current (ac) power supply with high frequency and high rates of gas flow recently has been subject of extended studies. This growing interest is justified by the fact that the APMJ are suitable for numerous applications that requires low gas temperature at glow regime of operation with low input power. Among them, one can emphasize the surface treatment of heat sensitive bio-medical materials, thin film deposition, reactive ion etching, bio-medical and dentistry applications, etc. For bio-medical applications, the non-thermal APMJ must be at glow discharge regime possibly at room temperature. It is desired to maintain the discharge at Townsend-glow regime to produce parallel-operated jets but the related current is very low, resulting in low power plasma which finally results in formation of a low density active species. The latter are essential for bacteriological decontamination since the formation of OH radicals and O atomsis favored. Kim et al proposed a microplasma device which operates at 20kHz ac power using a commercial transform for a neon light coupled with a set of ballast resistor and autotransformer in order to avoid excessive currents [1]. A flow of 10L.min⁻¹ nitrogen gas was injected through the microtube and ac voltage of 13.5kV_p-p was applied to produce the jet, conditions suitable for use in biomedicine, as a cancer therapy. Szili et al used microplasma jet to modify surface at different treatment conditions [2]. To power the microjet at atmospheric pressure, a power delivery system consisting of a sinusoidal oscillator and an autotransformer (10kV_p-p, 40mA_p-p, 40kHz) was used. When polystyrene surface is plasma treated by this method, the contact angle measurements reveal increase of the surface energy transforming it as hydrophilic. Ayan et al developed an atmospheric pressure non-thermal microjet for localized surface treatment [3]. The applied power ranged from 0.1 to 0.2W, useful for polyethylene substrate treatment without causing any damage. The gas mixture consists of 0.2L.min⁻¹ He + and 0.025L.min⁻¹ O₂ and surface functionalization (physical and chemical) was achieved. Seo et al designed a low-frequency (20 kHz, high voltage and mA current) He and Ar APMJ and measured its optical and electrical characteristics [4]. The high intensity of OH radical and O atoms in Ar plasma was suitable for large area sterilization and indirect treatment. Microplasma jet that operates at an excitation frequency of kilohertz ac range was developed to create nonthermal glow plasma. Here, He and O₂APMJ were investigated by adding O₂ flow, suitable for biomedical
applications. He non-thermal microplasma jets generated at 50 kHz ac voltage were electrical and optically characterizes [5]. It has been detected a number of highly reactive species in the plasma plume such as OH, O, N₂, He and N₂⁺ at the glow discharge mode. The presence of numerous of Helines indicates higher plasma density and/or electron temperature. Great disadvantage of these microplasma jets is the high cost of both the He gas (generally L.min⁻¹ He flow are needed to obtain a stable APMJ) and the equipment (high frequency power source used). The large surface to volume ratio of these miniaturized jets makes impossible surface treatments of a large scale. Furthermore, up to date there are not enough studies to detect and reducing the emission of a number of radicals such as HC, CO, NOₓ, O₃ and SOₓ, harmful to health if inhaled and manipulated without a due care.

Recently, an APMJ generated from micro-discharges in porous dielectric [6] or from the gap between two disk-shaped electrodes with a hole (1 mm of diameter) were presented [7]. Both used N₂ and air to generate glow plasma powered by a high-voltage ac power supply and a commercially available transformer for neon light operated at 60Hz connected to the electrodes. These gases have higher dielectric breakdown voltages compared to He or Ar and in static mode (without gas flow) it is practically impossible to produce a glow discharge at atmospheric pressure. After the breakdown transition occurs from a glow discharge to a filamentary glow or silent electric discharge [8, 9] in a very short time due to this electrodes configuration. To avoid thermal instability leaving to a filamentary glow or glow-to-arc transition (transition from glow to thermal discharges), they generate the non-thermal and stable glow discharge through large gas flow rate (5 – 15L.min⁻¹) that cools the plasma region removing the hot chemical species and neutrals. For surface treatment the presence of metastable states of some gases such as He (19.8 eV), Ar (11.6 eV) and N₂ (6.2 eV) is very interesting to produce microplasma. These energetic species react with the surface chemical structure with large efficiency. In this work, a 60 Hz ac APMJ device was studied to produce stable cold plasma suitable for surface treatment. The electrical and optical characteristics will be raised for a range of input power and for these gases.

2. Experimental

Figure 1 show the apparatus used to produce the ac microplasma jet. One transformer for neon light was used, operated at 60Hz. One can see in this figure the Cu electrodes, arranged parallel and inserted in a quartz tube. The nozzle cross-section area is about 5.3mm² and the gap between the circular electrodes is about 1.2mm. The ac power supply (12kV, 40mA) is coupled with a variable autotransformer and with a ballast resistor (50kΩ, 100W) in order to avoid excessive current and to allow voltage control. This figure presents the setup used to perform optical emission spectroscopy (OES) measurements of APMJ. The plasma was confined between the electrodes, burning freely in air at atmospheric pressure (709Torr in the laboratory and 46% local humidity) and was composed by the main gas and, in some cases, high-purity 1% H₂ was added. The addition of H₂ makes possible to estimate the electron number density of the plasma from the broadening of the hydrogen atom H β line. The gas connections were controlled by means of mass flow controllers MKS 247C and MKS model 1159B. The digital oscilloscope used in the experiments was a DPO7254 from Tektronix, 2.5 GHz. The highvoltage probe (from Tektronix) was a P6015A and a current probe TCP 312 with amplifier TCPA 300. The monochromator used to capture the light from the APMJ was a IHR550 model from Horiba-Jobin-Yvon with a diffraction grating of 1800 linesmm⁻¹ blazed at 500nm and 2400 linesmm⁻¹ at 300nm. The entrance slit was fixed at 120µm and the respective instrumental Gaussian-type full-width at half-maximum (FWHM) Δλᵢ = 0.1nm was obtained by means of an Ar-Hg low-pressure scientific lamp and analyzing the fitting procedure for the Hg 435.8 nm strong line broadening. The spectra at the exit slit were recorded by a CCD Symphony with 1024×512 pixels and the SynerJy software was used to handle the spectra.

Figure 1. Schematic representation of the microplasma jet and the schematic representation of the electrical circuit (the inset left represents the photograph of air plasma jet)
3. Results and Discussion

The frequency of the sinusoidal ac power supply is 60Hz and the glow discharge is achieved at a gas flow varying from 3 to 10 L.min⁻¹. In Fig. 2 a typical current-voltage waveforms of Ar APMJ is presented, at low rate of about 6 L.min⁻¹. The discharge voltage is about \( V_{pp} = 1200V \) (\( V_{rms} = 464V \)) for current of \( I_{rms} = 12.5mA \). A filamentary discharge is observed at gas flow below 2L.min⁻¹. The average electric power consumption was calculated by the equation \( P = \frac{1}{T} \int_{0}^{T} V(t)I(t) dt \) and results 4.8W.

Current and voltage measurements of stable APMJ at flow rate of 6 L.min⁻¹ show sinusoidal wave forms for flows of Ar, He, O₂ + 1%Ar, N₂ and air gases, with rmscurrent and voltage amplitudes of: a) 464 V, 10.4 mA; b) 410 V, 10.6mA; c) 515 V, 13 mA; d) 430 V, 11.8 mA and e) 456 V, 12.9 mA respectively. The phase shift is zero and the power absorbed by the plasma is: a) 4.8 W; b) 4.1 W; c) 5.8W; d) 4.4 W and e) 4.9 W respectively. The jet elongates axially when the flow rate is increased, from 3 to 10 L.min⁻¹. The highly reactive species as OH, O, N₂⁺ and energetic photons produced between the electrodes will extent along the plasma plume, both radially and axially from the exit nozzle of the APMJ.

Figure 3 shows the main atomic and molecular emission lines from the plasma when distinct gases are used individually or mixed. For individual mixtures between O₂, N₂ and air with 1% Ar and for pure He and Ar working gas (added with 2% O₂), the intense lines of Ar I 4p – 4s and 5p – 4s transitions (415.8, 416.4, 418.2, 419.1, 419.8, 420.1, 425.9, 427.2, 427.8, 433.4, 437.1, 437.6, 437.9, 440.0, 442.6, 444.9, 447.3, 452.2, 706.7, 714.7, 727.3, 738.4, 750.3, 751.5, 763.5, 774.8, 800.6, 801.5, 810.4, 811.5nm) and He emission spectrum (388.9, 396.5, 438.8, 447.2, 492.2, 501.6, 504.8, 587.6, 706.5, 728.1 and 776.8 nm) and Na I doublet lines at 589.59nm have been observed. One can see intense Ar I 4p – 4s transition lines from 700 to 800 nm, as well as the blue transitions 5p – 4s from 410 to 430 nm when pure Ar is used. From 420 to 450 nm, the Ar II emission lines represent 4p – 4s and 4p – 3d transitions originating from higher energy levels. The OH (first order, ultra-violet, \( A^3Σ^+ \), \( v = 0 → X^1Π, v = 0 \)) at 306.357 nm, NO, NH, O (616; 777 nm, 3p – 3s), N (750, 825 and 875nm), N₂ (\( C^0[1]u, v = 0 → B^2[1]g, v = 0, \epsilon = 1.1 eV, second positive system), N₂ (\( B^2[1]g → A^2Σ^+ \), first positive system), N₂⁺ (\( B^3[1]Σ^+ \), \( v = 0 → X^3Σ^+ \), \( v = 0 \)) at 437 to 475nm and H₂ (486.13nm) spectra and energetic UV photons were observed in the plasma plume. The reference [6] presents ac APMJ at 60Hz, similar to the others discussed above. They observed the later lines and O line at 844nm, this last representing a very reactive chemical species useful for bacteriological control. It is important to use a mixture of He or Ar gas with a little amount of O₂, because oxygen is required to produce active species aiming polymer surface treatment, especially at high pressure. Pure O₂ plasma at atmospheric pressure cause quenching of the discharge and this gas alone generate unstable plasma, with filamentary glow aspect. Therefore, the mixture of this gas with noble gases is important to produce stable APMJ, carrying selective oxygen species. Both the O atoms and OH radicals are interesting species for biomedical applications. In open-air atmospheric pressure discharges due to gas impurities and from the air wrapping up the plasma volume, most of these rotational lines are observed commonly. The analysis of the emission lines of some of these molecules and atoms, such as Ar and H, allows estimation of the most important microplasma parameters, such as modes temperatures along the plasma jet length for a range of currents. In Fig.3, the Na lines were observed when the sugarcane bagasse amount were pretreated with the plasma, using a gas mixture composed with air/Ar, aiming to convert cellulose into sugar (this application will be reported in a future work). The presence of Na I doublet lines at 588.99 and 589.59nm indicates collisional processes between the active species of the gas and the treated surface. These processes give an estimation of the energetic species reacting with large efficiency against the surface chemical structure of the sugarcane bagasse, because these lines represent the transition from the ground state to the upper level at 2.1eV. In the He APMJ driven by 4.8W power consumption and flow rate of 6L.min⁻¹, added with 2% O₂, we observed He I, O I and nitrogen atomic and molecular lines as N (750, 825 and 875nm), N₂ (\( C^0[1]u, v = 0 → B^2[1]g, v = 0 \)), N₂ (\( B^3[1]g → A^3Σ^+_u \)) and intense N₂⁺ (\( B^3[1]Σ^+ \), \( v = 0 → X^3Σ^+ \), \( v = 0 \)) at 391nm (see Fig.4).

Besides the He lines, the later lines appear as a result of the open air operation, where excited species of the plasma (electrons and metastables) interact chemically with the air gas neutrals, mainly N₂ and O₂. Ionic N₂⁺ at 391nm and atomic oxygen lines could be result from the following reactions [10]

\[ \text{He}^+ + \text{N}_2 \rightarrow \text{He} + \text{N}_2^+ + e^-, \text{Penning ionization} \]
\[ \text{He}^+ + \text{N}_2 \rightarrow \text{He} + \text{N}_2^+ , \text{charge transfer} \]
\[ e^- + \text{N}_2^+ \rightarrow \text{N}_2^{2+} + e^-, \text{direct electron impact excitation} \]
\[ \text{N}_2^+ + \text{O}_2 \rightarrow \text{N}_2 + 0 + 0, \text{Penning ionization} \]

When 2% of H₂ is added with the He flow, both the H₂ line intensity and FWHM are increased. H atoms may be formed by reactions involving ionic He in H₂ or in He. H⁺, H₂⁺ and H₂⁻ ions could be responsible for the observed Balmer atoms through the following reactions [11]

\[ \text{H}_2^+ + \text{H}_2 \rightarrow \text{H}_3^+ + \text{H} \]
\[ \text{H}_2^+ + \text{He} \rightarrow \text{H}_3^+ + \text{H} + \text{He} \]
\[ \text{H}_2^+ + \text{He} \rightarrow \text{He}^+ + \text{H}_2 \]

H₂ may be dissociated by collision with electrons and may produce H₂. The excitation of the H atoms should be done by electron impact

\[ H + e \rightarrow H^* + e \]

The spectra of OH bands were observed due to water vapor impurities in the gas and in the atmospheric air. The populations of excited state OH (\( A^2Σ^+, v = 0 \)) at 306.4 nm can be generated by direct dissociation electron excitation of water or by dissociative recombination of H₂O⁺.
where this radical is formed by He metastables and subsequent dissociative recombination [12]

\[
\begin{align*}
H_2O^+ + e^- &\rightarrow OH(A) + H + e^- \\
H_2O^+ + e^- &\rightarrow OH(A) + H \\
\end{align*}
\]

Figure 5 show the \( \text{N}_2^+ (B^2\Sigma_u^+, \nu = 0 \rightarrow \chi^2\Xi_g^+, \nu' = 0; 380 \text{ to } 475 \text{ nm}) \) for Ar and He flow (added with \( 2\% \) oxygen gas) at the same power consumption, 4.8W. One can see the \( \text{N}_2^+ \) \((B^2\Sigma_u^+, \nu = 0 \rightarrow \chi^2\Xi_g^+, \nu' = 0) \) at 391 nm is practically suppressed when the plasma jet is operated with \( \text{Ar} + 2\% \text{O}_2 \), because this APMJ spend more power to produce oxygen atoms, reducing the production of vibrational species like \( \text{N}_2^+ \) [10]. The Ar I 603.2 nm line broadening (FWHM of the Lorentzian shape of the line, \( \Delta \lambda_{\nu}\text{nm} = 3.9T_{g}^{0.7} \)) is strongly gas temperature dependent and the weakness of this line intensity is an indication of the low gas temperature of the microplasma jet.

![Figure 2](image-url)  
**Figure 2.** Voltage and current waveforms of the Ar microplasma jet at \( V_{rms} = 464 V \) and \( I_{rms} = 10.4mA \) and flow rate of \( 6 \text{ l/min} \)^{-1}

![Figure 3](image-url)  
**Figure 3.** OES measurements for a range of gas flow at APMJ
Figure 4. $N_2^+$ at 391nm for He and Ar microjet in open air

Figure 5. $N_2^+$ ($B^2Σ_u^+, \nu = 0 - X^2Σ_g^+, \nu' = 0$) at 391nm

$N_2^+$ ($B^2Σ_u^+, \nu = 0 - X^2Σ_g^+, \nu' = 0$) at 391nm for Ar and He flow at the same power consumption, 4.8W
The gas temperature $T_g$ was estimated from analysis of the OH (first order, ultra-violet, $A^2\Sigma^+\to X^2\Pi$, $\nu = 0\to0$) rotational band. Figure 6 presents the experimental rotational spectrum of OH (first order, ultra-violet, $A^2\Sigma^+\to X^2\Pi$, $\nu = 0\to0$) at 306.4 nm obtained for Ar APMJ. The $T_{\text{rot}}$ may be taken from the Boltzmann plot (BP) [13-19]. The spectra of OH bands were obtained with a resolution of 0.05 nm. These OH bands could be observed due to water vapor impurities in the gas, adjacent equipment and in the atmospheric air. The BP was used to obtain this molecule rotational temperature for $I_{\text{rms}} = 10.4$ mA, $T_{\text{rot}} = (448 \pm 33)$ K and flow of 6 L.min$^{-1}$ for low rotational levels energy $J$ ($Q_1$ branch). The APMJ gas temperature corresponds to the low OH rotational temperature. The gas temperature was kept constant due to the argon gas flowing between the electrodes and due to thermal conduction to the neighborhood. The plasma parameters are: gas flux $\Phi\sim 6$ l.min$^{-1} = 1.0\times10^{-4}$s$^{-1}$; gas density $N = P/k_B T_g = 1.49\times10^{25}$ m$^{-3}$; the number of argon atoms heated per second in the plasma is $n = \Phi N = 1.49\times10^{21}$ s$^{-1}$; the energy transfer to these atoms from 300 to gas temperature is $\Delta E = 0.5k_B\mu(8\pi\nu)^{(1/2)}(T_g - T_{\text{room}}) \sim 4.6$W, consistent with the discharge power. Figure 7 shows the second positive system ($C^3\Pi_u \to B^3\Pi_g$, $\nu = 0 \to \nu = 0$, $\epsilon = 11.1$ eV) of these molecule and Fig.8 the rotational spectrum of the second positive system of nitrogen (from 334 to 336 nm - R branch) applied to estimate the gas temperature in the plasma, where Ar metastables and excited species are present. For rotational lines belonging to the rotational levels $J$ varying from 14 to 28 of the R branch of the second positive system of nitrogen can be plotted 

$$\ln(I / J^{-1}) \propto B_v J (J + 1)h c / k_B T_{\text{rot}},$$

where the rotational constant $B_v = 1.815$ cm$^{-1}$. The linear fit of the Boltzmann plot gives the $N_2$ rotational temperature. Figure 9 depicts the measured $N_2$ rotational temperature obtained from Boltzmann plot at the same conditions of Fig.6, showing higher results for $T_{\text{rot}}$ by reasons explained before.

Figure 10 shows the gas temperature as a function of the discharge power, for a range of gases flowing through the microjet. The gas flow was maintained constant at 6L.min$^{-1}$. As stated before, this temperature was estimated from OH radical where we consider their rotational states and supposing that the neutral gas heavy particles are in equilibrium due to the low energies needed for excitation at this rotational level. One can see in this figure the increase rate of the temperature for all gases. For low input power the temperature is as low as the room temperature. He discharges at 5 to 6 L.min$^{-1}$ gas flow show good stability and low temperature, ideal for biomedical applications. Reference [6] refers to a $N_2$ feed ac APMJ operated at 60Hz. The authors found that the gas temperature varies from 300 to 500K at the nozzle exit. Optical emission spectroscopy measurements allow us to infer the mean gas temperature of plasma jet from the outer electrode along its axial position. A temperature decrease up to 40°C was observed up to 5 mm and beyond this the gas temperature reaches room temperature, so enabling the APMJ to use in treatment of heat-sensitive surfaces such as biomaterials and polymers. These last measurements will be reported in a future work. Figure 3 shows a set of neutral and ionic particles, especially highly reactive ones as OH, O and $N_2^+$. We can estimate the electronic temperature through Line Intensity Method, where the Saha equation is used to obtain the ratio between emitted spectral lines [13-17]. We choose the following lines (Fig.3): a) Ar$^+$ 4p-4s and Ar 4p-4s; b) $N_2^+$ (391nm) and $N_2$ rotational lines. Figure 11 shows the average electronic temperature which varies between 0.50 and 0.70eV for the full range of the gases flowing through the APMJ.

![Figure 6](http://example.com/Figure6.png)
Figure 7. Second positive system of N$_2$ from 300 to 460nm

Figure 8. Rotational spectrum of the N$_2$ ($C^3Π_u$, $\nu = 0 \rightarrow B^3Π_g$, $\nu' = 0$) transition (R branch)
Figure 9. Boltzmann Plot of the the N$_2$(0 – 0) ($C^3Π_u$, $\nu = 0 \rightarrow B^3Π_g$, $\nu = 0$) (R branch) rotational transitions (see Fig. 8)

Figure 10. Gas temperature as a function of the input power supplied to the discharge (6 L.min$^{-1}$ gas flow)
Figure 11. Electronic temperature of the discharge versus the electrical input power

Figure 12. Vibrational distribution of the $N_2 (C^3Π_u \rightarrow B^3Π_g)$
The Boltzmann plot method can also be applied to estimate the vibrational temperature ($T_v$) in the plasma bulk using the second positive system $N_2 (C^3Π_u \rightarrow B^3Π_g)$ for $\Deltaν = -2$. Figure 12 shows the vibrational transitions for this molecule with the corresponding band heads (input power $P = 8$ W). The temperature was in the range 0.35 to 0.57 K for the input power varying from 3 to 10W. Therefore, the difference between these three modes of temperature, $T_e>T_v>T_g$ indicates the partial-local thermodynamic equilibrium of the APMJ.

Figure 13 show the relative intensity of the H$_β$ line, Ar, air and He gases flowing for the same electrical conditions, namely at 4.8 W power consumption. Although there are distinct relative intensity for each flow, the FWHM of the Voigt shape line is $\Deltaλ_v = 0.12$ nm [13]. Unfortunately the Instrumental broadening, related to the opening slit of the monochromator, is approximately equal to 0.1 nm and it is impossible to obtain the gas temperature neither the electron number densities with sufficiently good accuracy using this method. Therefore, one can conclude from this figure that the gas temperature must be low ($T_g < 500 K$) and it is practically the same for the different gases. Compared to the dc glow discharge studied before [13], the electron number density is below than $10^{15}$ cm$^3$. According to the plasma parameters, $n_e = j(\varrho \mu_e)^{1/2} \approx 10^{12}$ cm$^3$.

4. Conclusions

The non-thermal low frequency (60 Hz) ac atmospheric pressure microplasma jet (APMJ) could produce stable cold glow discharges using numerous of gases such as Ar, He, N$_2$, O$_2$ and air. The APMJ is suitable for surface treatment. The jet has a visible radial diameter of approximately 1.5 mm with afterglow surrounding the jet. Current and voltage measurements at flow rate of 6 L.min$^{-1}$ show sinusoidal wave forms for flows of Ar, He, O$_2$ + 1%Ar, N$_2$ and air gases, with rms current and voltage amplitudes of: a) 464 V, 10.4 mA; b) 410 V, 10.6 mA; c) 515 V, 13 mA; d) 430 V, 11.8 mA and e) 456 V, 12.9 mA respectively. The power absorbed by the plasma was: a) 4.8 W; b) 4.1 W; c) 5.8 W; d) 4.4 W and e) 4.9 W, respectively for the mentioned species. The optical measurements show the emission of important lines for bio-medical and surface treatment applications, as O, N$_2$, O$_3$, N$_2^+$, N and OH. Electronic temperature ($T_e$) was estimated through the Line Intensity Method. The average electronic temperature varies between 0.50 and 0.70 eV for a vast ranges of gas flow. The Boltzmann plot method was applied to estimate the vibrational temperature ($T_v$) in the plasma bulk using the second positive system $N_2 (C^3Π_u \rightarrow B^3Π_g)$ for $\Deltaν = -2$. The temperature was in the range 0.35 to 0.57 K for the input power varying from 3 to 10W, for all gases flow. The gas temperature $T_g$ was estimated analyzing the OH (first order, ultra-violet, $A^2Σ^+, ν = 0 → X^2Π, ν = 0$) rotational band. A temperature decrease up to 40°C was observed up to 5 mm and beyond this range the gas temperature reaches the room temperature, enabling so the use of APMJ in treatment of heat-sensitive surfaces such as biomaterials and polymers. The difference between these three modes of temperature, $T_e>T_v>T_g$, demonstrates near PLTE condition of this plasma source. The capacity of operation with molecular and atomic (noble) gases, including air, is important for various technological applications, but the novelty of this APMJ is
the broadening of the plasma plume at the outer electrode with diameter of 2.5 mm and, consequently, reaching higher surface treatment area.

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