Doped TiO₂: the effect of doping elements on photocatalytic activity†

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Doping of TiO₂ with various elements increases its photocatalytic activity due to the formation of new energy levels near the conduction band. Photocatalysis involving titanium dioxide is a heterogeneous process in which the surface of the catalyst plays an important role. The structural properties of TiO₂ are influenced by the synthesis method, the doping method, and the dopants. In this work, we compare different doping elements for improving the photocatalytic activity of titanium dioxide, which was synthesized by the sol–gel method. In the doping method, low-temperature underwater plasma was used. Al, Cu, Mo, and W acting as electrodes were chosen as doping elements. The obtained samples were characterized by various techniques. The incorporation of elements leads to the distortion of the TiO₂ crystal lattice, thus changing its surface characteristics, and to a decrease in the band gap. The introduction of aluminum and copper increases the photocatalytic activity to 70% while doping with Mo and W increases the activity to 96% upon visible light irradiation for 60 minutes. Explanations of the effect of various doping elements on the photocatalytic activity of titanium dioxide are presented.

Introduction

In the processes of heterogeneous catalysis, surface area plays the main role. In the processes of photocatalysis, the redox potential of charge carriers and the selectivity of the catalyst surface are important. Recently, nanostructured materials have been considered for these purposes. The most popular of these is titanium dioxide.1,2 TiO₂ is a semiconductor whose photocatalytic properties deteriorate under visible light due to its wide band gap. It is known that doping increases the photocatalytic activity of TiO₂ upon irradiation with visible light. The nature of the dopant ion also affects this activity since the dopants can be interstitial, substitutional, or both in the case of incorporation. Different locations have a different effect on the properties of titanium dioxide. A comparison of the effectiveness of photocatalytic decomposition or synthesis of organic compounds is difficult since doping is carried out using different methods (sol–gel, solid reaction/mechanical activation, chemical vapor deposition, etc.), in which various precursors of metal ions are used.3–7 This can also affect the photocatalytic activity of the resulting samples.

TiO₂ is doped with various elements, both metals and non-metals.8,9 The most popular are transition metal ions.10–14 Interest in such elements is caused by the fact that they have partially filled d-orbitals. The incorporation of them into the lattice of titanium dioxide induces the formation of new energy levels near the conduction band. The use of Cu is a more affordable and cheaper alternative to elements such as silver and gold, as well as platinum group metals.15–19 Recently, there has been increasing interest in the use of trivalent metals (Al, Nb) as doping materials as they improve the electrical, optical, and structural characteristics of titanium dioxide.20–24 Doping of Al, Cu, Mo, or W always leads to a narrowing of the band gap of TiO₂, which increases its photocatalytic activity upon irradiation with visible light. The introduction of Al³⁺ into the crystal lattice of titanium dioxide leads to the appearance of oxygen vacancies, which increase the photocatalytic activity.25 It was shown that, upon doping with copper, clusters of copper oxide are formed on the surface of TiO₂, which can also take part in the photocatalysis.16,17 In the cases of transition metals of molybdenum and tungsten, not only Mo⁶⁺ and W⁶⁺ ions but also Mo⁵⁺, Mo⁷⁺, W⁴⁺, W⁵⁺, as well as Ti⁷⁺ ions, which also take part in photocatalytic processes, can be present on the powder surface.26–29

In the present work, we report a green chemistry method to obtain metal-doped titanium dioxide by the one-step plasma treatment of the TiO₂ sol. The main goal was to study the effect of the doping cation on the structural and photocatalytic properties of titanium dioxide. TiO₂ was prepared using a low-temperature sol–gel method, which allows one to obtain micro- and mesoporous powders. The doping procedure was...
carried out using low-temperature underwater plasma with metal electrodes. Al, Cu, Mo, and W were chosen as the sources of doping elements. The obtained samples were characterized by different techniques. The effect of the metal ion doping on the visible light photocatalytic performance of doped TiO$_2$ was estimated towards the Rhodamine B dye.

The novelty of the work is the use of a new method for doping titanium dioxide with various elements by the plasma exposure of the TiO$_2$ sol. Usually, the synthesis of titanium dioxide and its doping are carried out by similar methods, for example, sol–gel synthesis and chemical doping. In this work, we propose a combination of sol–gel synthesis with plasma treatment, during which the TiO$_2$ sol is modified and doped. We use simple and affordable materials as doping elements. In this method, no chemical reagents are used as metal precursors. It increases the purity of the samples. There is no need to anneal the sample for cleaning when using this method. As a strength of this work, it can be noted that the effect of the nature of the doping element and the charge state of doping elements. The obtained samples were characterized for investigating the effect of the initial concentrations of dyes and contact time while studying the adsorption isotherm and kinetics of the reactions. Dye solutions with known concentrations were prepared. The concentration of the dye was 1.2–12 mg L$^{-1}$. The required amounts of powders were added to the flasks with solutions. The dispersed solutions

**Materials and methods**

**Materials**

Al, Cu, Mo, and W wires (diameter 1 mm), titanium isopropoxide, and Rhodamine B were purchased from Aldrich (USA). Nitric acid and isopropanol were obtained from Chimmed (Russia). All chemicals were used as received.

**Methods**

The synthesis of TiO$_2$ was carried out using a sol–gel method. 16 mL Ti isopropoxide (98%) was mixed with 12 mL isopropanol for 2 hours with a magnetic stirrer. The next stage was adding an aqueous solution of nitric acid at pH 2. The mixing was carried out for 4 hours at 80 $^\circ$C for sol and gel formation. The prepared TiO$_2$ sol was processed by plasma. The time of treatment was 3 minutes. All experiments were carried out with a 100 mL volume of liquid. Then the suspension was dried at 60 $^\circ$C in an oven for 24 hours. A scheme of the synthesis and doping of TiO$_2$ is presented in Fig. 1.

Underwater plasma is excited between two metallic electrodes, which were immersed in the liquid sol. Al, Cu, Mo, and W wires with a diameter of 1 mm were used as metallic electrodes. The distance between electrodes was constant (3 mm). The plasma was ignited using a direct current (DC) power supply. The masses of the electrodes before and after the experiment were measured using an AND HR-15AZ analytical balance (error $\pm$0.1 mg).

**Characteristics of the synthesized samples**

The suspension solution obtained during the action of the discharge was analyzed by dynamic light scattering to determine the size of the formed particles and the zeta potential using a Zetasizer Nano (Malvern). The absorption spectrum was recorded in the UV and visible spectral regions to estimate the band gap.

X-ray diffraction experiments were carried out on a diffractometer with a Mo-K$_\alpha$ source ($\lambda = 0.07107$ nm) (Bruker D8 Advance). The radiation was monochromatized by a zirconium $\beta$ filter. Diffraction data were recorded in the range of angles 4–40 $^\circ$ (2$\theta$) with a step of 0.02 $^\circ$. The intensity of scattering from the samples was registered using a VANTEC-1 position-sensitive detector (Bruker). The exposure time at each experimental point was 0.5 s. The crystallite sizes were estimated using the Debye–Scherer equation (eqn (1)):

$$D = \frac{k\lambda}{\beta\cos \theta},$$

where $\lambda$ is the wavelength of the X-ray radiation (0.07107 nm), $\beta$ is the full-width at half maximum after the subtraction of instrumental noise (free of instrument induced broadening), and $\theta$ is the diffraction angle. An adequate baseline correction was used to obtain the best values of $\beta$ (and as a consequence $D$). The surface of the crystals and their elemental composition were examined using a scanning electron microscope (TESCAN VEGA 3 SBH, Czech Republic) equipped with an energy dispersive spectrometer (EDS) (Oxford Instruments, UK). The specific surface area and pore size were determined using low-temperature nitrogen adsorption–desorption (NOVA Series 1200e).

**Adsorption and photocatalytic tests**

Rhodamine B adsorption studies were carried out in batch form for investigating the effect of the initial concentrations of dyes and contact time while studying the adsorption isotherm and kinetics of the reactions. Dye solutions with known concentrations were prepared. The concentration of the dye was 1.2–12 mg L$^{-1}$. The required amounts of powders were added to the flasks with solutions. The dispersed solutions
were mixed using a magnetic stirrer at 300 rpm. Then the solutions were filtered and the residual concentration of RhB was determined using a UV-Vis spectrophotometer (PG70+ Instruments) at 544 nm. The adsorption kinetics was determined by analyzing the adsorption capacity of the dye solution at different contact times.

The photocatalytic activity of doped TiO₂ was evaluated through the rate of destruction of the Rhodamine B (RhB) dye under visible light. The experimental setup for the photocatalytic destruction of the dye has been described in detail.³¹

The source of visible light irradiation was a 55 W lamp (Phillips). 0.03 g of the powder was added into 500 mL of the 1.2 mg L⁻¹ RhB solution. The dye solution and powder were mixed using a magnetic stirrer.

### Results and discussion

#### Characteristics of the powders

Data of the X-ray diffraction analysis are presented in Fig. 2. The interpretation of data was done by using the crystallographic open database.³² The untreated TiO₂ powder showed characteristic peaks related to the (101), (112), (200), (211), and (204) planes at 2θ = 11.5°, 17.36°, 21.54°, 24.93°, and 27.43°, respectively. All peaks correspond to the anatase phase of TiO₂. No additional peaks related to impure metals or their oxides have been recorded. A little shift of the main diffraction peaks was detected in powders after plasma treatment. This means that the doping of titanium dioxide occurs during plasma treatment. The lattice parameters are calculated using eqn (2) for the (200) and (101) peak positions,³³

\[
\frac{1}{d^2} = \frac{h^2}{a^2} + \frac{k^2}{b^2} + \frac{l^2}{c^2}
\]

where \(h, k, l\) are Miller’s indexes and \(d\) is an interplanar distance. The calculated data are presented in Table 1.

| Sample          | Lattice parameters, Å | Crystallite size, nm |
|-----------------|-----------------------|----------------------|
| TiO₂            | 3.78 9.50 135.74      | 4.48                 |
| Al/TiO₂         | 3.79 9.46 135.88      | 4.66                 |
| Cu/TiO₂         | 3.77 9.39 133.46      | 4.61                 |
| Mo/TiO₂         | 3.78 9.50 135.74      | 4.13                 |
| W/TiO₂          | 3.80 9.46 136.60      | 3.89                 |

Table 1  Lattice parameters and crystallite size of the synthesized samples

See that the cell parameters of untreated and treated TiO₂ are close but the \(a\) and \(c\) parameters change slightly, which results in cell volume changes. The changes in the values of the cell volume lie within the limits of calculation errors (1%). The ionic radii of Mo⁶⁺ (0.59 Å) and W⁶⁺ (0.60 Å) are close to that of Ti⁴⁺ (0.605 Å). Therefore, substitution with these elements should not lead to changes in the unit cell. The ionic radius of Al³⁺ is smaller (0.535 Å) than that of Ti⁴⁺. For Al/TiO₂ the lattice parameter \(a\) is increased and the parameter \(c\) is decreased as reported by Xu et al.³⁰ The ionic radius of Cu²⁺ (0.73 Å) is slightly larger than that of Ti⁴⁺. It was shown in ref. 34 by the DFT method that the introduction of a copper ion into the TiO₂ lattice leads to a decrease in the linear distance between oxygen atoms. If we assume that the introduction of the copper ion occurs in the internode, which is the center of compression, we can explain the decrease in the unit cell.

Analysis of the published data showed that while doping with Mo⁶⁺ and W⁶⁺ an increase in cell volume and a decrease in crystallite size are always detected. In the case of aluminum as a doping agent, the published results are not so clear. In ref. 14 and 21, it was established that the introduction of Al led to an increase in the crystal lattice parameters \(a\) and \(c\), while in ref. 24 the opposite effect was found. Published data on aluminum doped with TiO₂ are contradictory. The reason for this may be the different sources of aluminum (salts, metallic aluminum).

In addition to the ionic radii of the substituting elements, their charges must also be taken into account. The loading of Al³⁺ in the TiO₂ lattice induces the substitution of Ti⁴⁺ and the appearance of an imbalance of charges. Oxygen vacancies are formed for the compensation of charge imbalance. They also cause lattice defects (tensile, compressive strain, etc.).²¹,³⁵ The introduction of W⁶⁺ or Mo⁶⁺ also leads to an imbalance of charges. In this case, an excess charge attracts hydroxyl ions to the surface to neutralize the charge.³⁶

In our case, the introduction of tungsten caused asymmetric changes in the lattice parameters, which could be the result of the substitution of Ti⁴⁺ and compensation of the charge imbalance. In the case of Mo/TiO₂, the calculations showed that there were no changes in the lattice parameters. This may be explained by the following. The underwater plasma action leads to the formation not only Mo⁶⁺ but also Mo⁴⁺ and molybdenum in the intermediate oxidation state.³⁷ The ionic radius of Mo⁴⁺ is close to that of Ti⁴⁺.³⁸ This can take the place of titanium inside the crystal lattice without causing a charge imbalance. A more detailed answer would probably be provided through...
XPS analysis. However, at present, we do not have the technical ability to conduct such an analysis.

SEM images with element mapping of doped TiO$_2$ are presented in Fig. 3. The distribution of doping elements on the titanium dioxide surface is non-uniform. The energy dispersive analysis showed that the content of the doping element did not exceed 1% (Table 2).

During the experiments, we believe that all the material of the electrodes atomized during the burning of the plasma was built into the lattice of titanium dioxide, thereby resulting in its doping. Based on the change in the mass of the electrodes before and after the plasma treatment of the sol ($\Delta m$), we calculated the maximum possible metal content (theoretical yields) in TiO$_2$ (Table 2). In experiments with aluminum electrodes, the value of the theoretical yield turned out to be lower than the experimental value. Taking into account the fact that the EDS analysis gives an error of about 2–3%, it can be argued that the theoretical and experimental yields coincide. The same can be said for experiments with the Cu and Mo electrodes. In the case of tungsten electrodes, it can be assumed that the theoretical excess of the atomized metal served as a crystallization center during the plasma treatment of the sol.

The isotherms of low-temperature nitrogen adsorption–desorption for all samples are of type I (Fig. S1 in the ESI†). The presence of hysteresis indicates the presence of micropores in the samples. This is also confirmed by the significant difference in the $S_{BET}$ and $S_{BJH}$ values (Table 3). The resulting powders have a high specific surface area. Except for the data for Cu/TiO$_2$, the introduction of a doping element increases the specific surface area of the powders. All doped materials are characterized by the presence of micropores (1.5–2 nm) and mesopores (3–4 nm) (Fig. S1 insets, ESI†).

In many studies, it was noted that the doping of TiO$_2$ leads to an increase in the specific surface of the obtained powders. In several works related to studies of the structures of titanium dioxide doped with copper, it was found that the introduction of Cu leads to a decrease in $S_{BET}$.$^{16,19}$ This is associated with the formation of copper oxide clusters that block pores.$^{16,39,40}$ The introduction of a doping element is always accompanied by the appearance of various defects in the crystal structure. The accumulation of defects (for example, oxygen vacancies) forms pores, thereby increasing the surface area.$^{41}$

### Table 2 Results of the EDS analysis and the theoretical yield of doping elements

| Sample    | Element ratio, % | Theoretical yield |
|-----------|------------------|--------------------|
|           | O     | Ti   | Metal | $\Delta m$, g | Yield, % |
| Al/TiO$_2$ | 48.28 | 51.48 | 0.24  | 0.009 | 0.21 |
| Cu/TiO$_2$ | 56.14 | 43.74 | 0.11  | 0.005 | 0.12 |
| Mo/TiO$_2$ | 48.56 | 50.99 | 0.46  | 0.02  | 0.54 |
| W/TiO$_2$  | 49.51 | 50.36 | 0.12  | 0.024 | 0.57 |

### Table 3 Structural characteristics of undoped and doped TiO$_2$

| Sample    | $S_{BET}$, m$^2$ g$^{-1}$ | $S_{BJH}$, m$^2$ g$^{-1}$ | $D_{av}$, nm | $V_p$, cm$^3$ g$^{-1}$ | $S_{mic}$, m$^2$ g$^{-1}$ | $V_{total}$, cm$^3$ g$^{-1}$ |
|-----------|-------------------------|--------------------------|--------------|------------------------|--------------------------|---------------------------|
| TiO$_2$   | 126.3                   | 18.1                     | 3.7          | 0.017                  | 26.1                     | 203.7                     | 0.069                  |
| Al/TiO$_2$| 128.7                   | 20.7                     | 3.6          | 0.019                  | 28.5                     | 210.8                     | 0.071                  |
| Cu/TiO$_2$| 105.5                   | 17.6                     | 3.2          | 0.016                  | 25.4                     | 176.3                     | 0.059                  |
| Mo/TiO$_2$| 137.9                   | 24.7                     | 3.6          | 0.021                  | 34.4                     | 239.9                     | 0.085                  |
| W/TiO$_2$ | 142.9                   | 20.7                     | 3.6          | 0.019                  | 29.9                     | 231.9                     | 0.078                  |

**Characterization of the liquid sol**

Fig. 4 shows the absorbance spectra of untreated and plasma treated TiO$_2$ sols. As it can be seen, the introduction of the
doping element leads to a red shift of the spectra. The relation (3) allows an estimation of the band gap energy (Table 4). The value of the band gap for TiO\textsubscript{2} was 3.28 eV. This value is in good agreement with the data presented in previous works,\textsuperscript{26,42} where it was found that the $E_g$ value is determined by the synthesis method.

The obtained values are in agreement with the published data for low concentrations of doping elements.\textsuperscript{22,24,27,42,43}

$$E_g = \frac{1240}{\lambda}.$$  \hspace{1cm} (3)

The dynamic light scattering method was used for the estimation of the average particle size and zeta potential in the colloidal media (Table 4). Particle size distribution histograms are given in the ESI\textsuperscript{*} (Fig. S2 in the ESI\textsuperscript{*}). It is shown that element loading leads to an increasing average size of particles and decreases the value of zeta potential.

The average particle size in the sol was determined by dynamic light scattering immediately after the plasma doping process. The basis of this method is the calculation of the hydrodynamic radius (size) of a particle by measuring the viscosity of the medium and its diffusion coefficient. According to the Stokes–Einstein equation, the particle size is inversely proportional to the viscosity of the medium. The underwater plasma treatment leads to a change/decrease in the viscosity of the sol by the partial destruction of the organic component of the sol. From this point of view, a slight increase in the particle size can be explained. The greatest effect was recorded for Mo/ TiO\textsubscript{2}. In this case, it can be assumed that molybdenum electrodes or molybdenum ions/atoms acted as a catalyst for the destruction of organic components.

**Sorption**

The sorption capacity of TiO\textsubscript{2} and doped titanium dioxide with Rhodamine B is depicted in Fig. 5. The quantities of the adsorbed dye were calculated using eqn (4). The sorption capacity of Mo and Al exceeds the capacity of undoped TiO\textsubscript{2} only. Unfortunately, we could not find the published data concerning the sorption capacity of doped TiO\textsubscript{2} towards organic dyes.

$$q_t = \frac{(C_0 - C_t) \cdot V_0}{m_{\text{cat}}},$$  \hspace{1cm} (4)

where $C_0$ and $C_t$ are the concentrations of the dye initially and at a predetermined time point, respectively (mg L\textsuperscript{-1}), $m_{\text{cat}}$ is the mass of the catalyst (0.03 g), and $V_0$ is the volume of the solution of the dye (0.05 L).

To analyze the sorption processes, we used models of pseudo-first and pseudo-second orders, the linear forms of the equations of which are of the form:

$$\ln(q_e - q_t) = ln\frac{q_e}{q_t} - k_1t,$$  \hspace{1cm} (5)

$$\frac{1}{q_t} = \frac{1}{k_2q_e} + \frac{t}{q_e},$$  \hspace{1cm} (6)

where $q_e$ and $q_t$ are the quantities of the adsorbed dye at equilibrium and at time $t$, respectively, $t$ is the contact time, and $k_1$ and $k_2$ are the constant rates of the pseudo-first and pseudo-second orders, respectively. Also, an intraparticle diffuse kinetic model was used, the linear form of the equation of which has the form:

$$q_t = K_{id}t^{1/2} + A,$$  \hspace{1cm} (7)

where $K_{id}$ is the diffusion coefficient and $A$ is the intraparticle diffusion constant that is directly proportional to the thickness of the boundary layer.

The pseudo-first ($k_1$ and $q_e$) and pseudo-second ($q_e$ and $k_2$) parameters were calculated from the plots of $\ln(q_e - q_t) = f(t)$

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Table 4  Properties of species in the liquid phase

| Sample   | $\lambda$, nm | $E_g$, eV | $D_{av}$, nm | $\zeta$, mV |
|----------|---------------|-----------|--------------|-------------|
| TiO\textsubscript{2} | 377.91 | 3.28 | 40.51 | 36.15 |
| Al/TiO\textsubscript{2} | 416.42 | 2.98 | 40.86 | 31.80 |
| Cu/TiO\textsubscript{2} | 400.11 | 3.10 | 40.95 | 31.83 |
| Mo/TiO\textsubscript{2} | 402.92 | 3.08 | 43.25 | 31.95 |
| W/TiO\textsubscript{2} | 387.89 | 3.19 | 41.03 | 31.85 |

**Fig. 4** Absorbance spectra of untreated and plasma treated TiO\textsubscript{2} sols.

**Fig. 5** Kinetic curves of RhB adsorption on TiO\textsubscript{2} and doped TiO\textsubscript{2} at room temperature.
The adsorption isotherms are shown in Fig. 6. For the Al/TiO₂ sample, the isotherm is of S type. For the Cu/TiO₂ and W/TiO₂ samples, the isotherms are of L type. In the case of Mo/TiO₂, the curve has the form of the initial part of the S isotherm.

The linear forms of the Langmuir and Freundlich model equations are as follows:

$$\frac{C}{q_e} = \frac{1}{K_L q_{\text{max}}} + \frac{1}{q_{\text{max}}} C,$$

(8)

$$\ln(q_e) = \ln(k_f) + nR \ln(C_e),$$

(9)

Table 5 Kinetics parameters and correlation coefficients ($R^2$) of RhB sorption onto the synthesized samples

| Sample   | TiO₂ | Al/TiO₂ | Cu/TiO₂ | Mo/TiO₂ | W/TiO₂ |
|----------|------|---------|---------|---------|---------|
| $q_e$ (exp.), mg g⁻¹ | 3.22 | 4.09 | 3.03 | 10.27 | 2.79 |
| Pseudo-first order | | | | | |
| $q_e$ (calc.), mg g⁻¹ | 2.48 | 4.91 | 3.12 | 6.6 | 2.74 |
| $k_1$, min⁻¹ | 0.445 | 0.661 | 0.467 | 0.611 | 0.502 |
| $R^2$ | 0.926 | 0.872 | 0.909 | 0.95 | 0.95 |
| Pseudo-second order | | | | | |
| $q_e$ (calc.), mg g⁻¹ | 3.43 | 4.88 | 3.59 | 11.01 | 3.41 |
| $k_2$, g mg⁻¹ min⁻¹ | 0.394 | 0.164 | 0.159 | 0.48 | 0.158 |
| $k_3$, g mg⁻¹ min⁻¹ | 0.995 | 0.983 | 0.992 | 0.992 | 0.99 |
| $R^2$ | 0.561 | 0.212 | 0.118 | 1.048 | 0.127 |

The calculated values of the constant and correlation coefficients ($R^2$) are presented in Table 5. The calculated values showed that the sorption process is not described internally by partial diffusion, since one of the three conditions is not satisfied. Because $R^2$ is close to 1 and the proximity of the $q_e$ (calc.) and $q_e$ (exp.) values, the sorption process is described using the kinetic model of the pseudo-second order.

Adsortion isotherms

The photocatalytic destruction experiments of the RhB dye solution with undoped and doped TiO₂ under visible light irradiation were carried out. The experiments were carried out in the dark and under exposure to visible light (Fig. 7). The results showed that the sorption process is not limiting. As can be seen, the introduction of a doping element in titanium dioxide leads to an improvement in its photocatalytic properties. In the cases of Al and Cu, the decomposition efficiency increases from 50% for pure TiO₂ to 70%. The loading of Mo and W increases the efficiency of photocatalytic destruction of the dye to 96% with 60 minutes of exposure.

$$q_e = \frac{C_e}{1 + \frac{C_e}{q_{\text{max}}}},$$

(Fig. S3 in the ESI†) and $t/\tilde{q}_e = \tilde{f}(t)$ (Fig. S3 in the ESI†), respectively. The parameters for the intra-particle diffusion model were obtained from the plot of $q_e = \tilde{f}(t)$ (see Fig. S4, ESI†). The approximation data are presented in Table 6. According to the values of the correlation coefficient for most samples, adsorption is not monolayer. The Freundlich isotherm model seems better than the Langmuir model for most samples. The $n$ values are more than 1. This means that the adsorption of one molecule raises the sorption of the other molecules of the dye. This is correlated with the properties of Rhodamine B, which can form dimers in low viscosity solvents (water) (so-called H-dimers).45

Photocatalytic activity

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$$k_f = k_2 C_e^{1/n},$$

where $k_f$ is the Langmuir constant related to sorption energy (L mg⁻¹), $q_{\text{max}}$ is the maximum capacity of the monolayer (mg g⁻¹), $k_f$ is the Freundlich constant which is connected with an affinity of the sorbed substance (ion) to the surface of the adsorbent (mg g⁻¹), and $n$ is the parameter of adsorption intensity.

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Photocatalytic activity

The photocatalytic destruction experiments of the RhB dye solution with undoped and doped TiO₂ under visible light irradiation were carried out. The experiments were carried out in the dark and under exposure to visible light (Fig. 7). The results showed that the sorption process is not limiting. As can be seen, the introduction of a doping element in titanium dioxide leads to an improvement in its photocatalytic properties. In the cases of Al and Cu, the decomposition efficiency increases from 50% for pure TiO₂ to 70%. The loading of Mo and W increases the efficiency of photocatalytic destruction of the dye to 96% with 60 minutes of exposure.
Registration of absorption spectra of dye solutions in the range of 190–800 nm during photocatalysis provides information not only about the destruction of the color centers of the dye molecule but also about the appearance of degradation products in the solution. Small fractions such as phenols and nitrophenols absorb in the region of 250–350 nm. In our experiments, the appearance of peaks in this region is not recorded (Fig. S5 in the ESI†). This allows us to assume that the destruction process leads to the formation of small organic molecules (saturated and unsaturated carboxylic acids). Data on the chemical analysis of solutions after photocatalysis can give more information about the destruction process. This was not the aim of this work, but may be the goal of future studies.

For a simplified description of the kinetics of heterogeneous processes, we make some assumptions. We accept the condition that the distribution of light intensity is the same throughout the reactor. Since we conduct experiments with one radiation source, the radiation intensity is the same in all experiments. The concentration of the catalyst is constant. Also, we will not take into account the rate of generation of active centers on the surface of the catalyst, which is determined by the volumetric absorption rate of photons. The photochemical reactor provides for cooling. Therefore, we accept the condition that the reaction temperature does not change and is 23 °C. Under these assumptions, the decomposition rate is determined only by the concentration of the substrate. The time dependences of \( \ln(C/C_0) \) are shown in Fig. 8. The kinetic curves are described using the pseudo-first-order reaction with good correlation coefficients \( (R^2 > 0.95) \). The slopes of the present lines show that the rate constant of W/TiO₂ is larger than those of other samples. The kinetic constants of undoped and doped TiO₂ are listed in Table 7.

The substitution of the Ti⁴⁺ ion or the incorporation of metal ions into the crystal lattice of titanium dioxide creates additional energy levels in the band gap of TiO₂. This causes photoactive transitions upon irradiation with visible light due to the excitation of electrons in this energy level and their transition to the conduction band. The use of metal ions as doping elements affects the photoreactive ability of titanium dioxide by acting as electron (holes) traps and changing the rate of the electron/hole recombination process. Electrons can be trapped by oxygen molecules on the surface of TiO₂ to form a radical O₂⁻. Holes move in the opposite direction from electrons. Photogenerated holes in the valence band can be scavenged by either hydroxyl ions or water molecules that are adsorbed on the surface. As a result, active OH radicals are formed, which are strong oxidizing agents.

According to the calculated values of \( E_g \), the effectiveness of doped titanium dioxide should decrease in the order of Al/TiO₂ > Mo/TiO₂ > Cu/TiO₂ > W/TiO₂ > TiO₂. However, we did not observe such a pattern. Therefore, the band gap is not the main criterion in our experiments. The specific surface plays a major role in sorption processes, which are not the dominant processes in our photocatalysis experiments.

Data from an XPS analysis can help us explain this phenomenon. Unfortunately, we do not have the technical ability to conduct such analysis of our samples. So we analyze the data presented in the published works. When doping titanium dioxide with molybdenum and tungsten ions, an XPS analysis shows that Mo⁶⁺, Mo⁵⁺, (in some cases Mo⁴⁺), W⁶⁺ and W⁵⁺ ions coexist on the surface. In addition to this, XPS peaks related to Ti⁴⁺ are recorded. All this leads to the formation of oxygen vacancies. In experiments on doping with aluminum ions, Al³⁺ and Ti⁴⁺ ions exist on the sample surface.

On the other hand, Aviles Garcia et al. found that the surface of the TiO₂ samples doped with molybdenum and tungsten contains surface hydroxyl groups. These groups are involved in the photodegradation of organic compounds. Also, in ref. 48 it was established that the surface acidity of W-doped TiO₂ is 42 times higher than that of Al-doped TiO₂. This explains the high affinity of the W/TiO₂ sample for chemical particles that have unpaired electrons.

Taking into account the fact that the state of the catalyst surface also plays an important role in heterogeneous photocatalysis, the following can be assumed. The action of plasma on aqueous and organic solutions leads to the formation of hydrogen peroxide. It is known that aluminum is inert concerning H₂O₂, copper is inactive, and Mo and W actively react in hydrogen peroxide with the formation of peroxy compounds and hydroxyl radicals. Based on this, we can conclude that the underwater plasma processing of sols leads not only to the doping of TiO₂ but also to the creation of new functional groups that take part in photocatalysis. In the case of using electrodes from Mo and W, this effect is most pronounced.

### Table 7: Rate constants \( (k) \) and correlation coefficients \( (R^2) \) of undoped and metal doped TiO₂

|          | TiO₂ | Al/TiO₂ | Cu/TiO₂ | Mo/TiO₂ | W/TiO₂ |
|----------|------|---------|---------|---------|---------|
| \( k, \text{min}^{-1} \) | 0.0092 | 0.0157 | 0.0169 | 0.0183 | 0.0351 |
| \( R^2 \) | 0.991 | 0.989 | 0.995 | 0.96 | 0.972 |

![Fig. 8: Kinetic curves of RhB destruction over undoped and doped TiO₂ under visible light irradiation.](image-url)
It is known that the time of diffusion of charge carriers from the bulk to the surface of the photocatalyst is directly proportional to the squared radius of particles.\textsuperscript{14,15,19,22,23,27,29,43,46,49} It means that a large crystallite size leads to an increase in the diffusion time. The probability of the electron/hole recombination increases, which leads to a decrease in the photocatalytic activity. A comparison of the data in Tables 1 and 6 shows that the photocatalytic activity of the samples decreases with increasing crystallite size.

It was noted that the effect of the doping amount on the photocatalytic activity has a curve with a maximum.\textsuperscript{14,15,19,22,23,27,29,43,46,49} The analysis of published works showed that the value of the optimal doping amount differs for various elements. In the case of Al/TiO\textsubscript{2}, the optimum amount can be 0.5 wt%\textsuperscript{22} or 0.2 wt%.\textsuperscript{23} These differences can be explained by both synthesis/doping methods and photocatalytic objects.

Conclusions

The use of underwater plasma as a method for doping titanium dioxide with various metal ions has shown good results. The photocatalytic ability was studied on the destruction of the Rhodamine B dye. All doped samples showed higher efficiency than undoped TiO\textsubscript{2}. The obtained results showed that the main role is played by the state of the catalyst surface in the process of photocatalysis. The surface is modified by the action of underwater plasma. Modification occurs in different ways and depends on the material of the electrode, which may be one of the directions of our future studies.

Conflicts of interest

There are no conflicts to declare.

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