Phase behavior of heavy oil–solvent mixture systems under reservoir conditions

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Abstract
A novel experimental procedure was proposed to investigate the phase behavior of a solvent mixture (SM) (64 mol% CH4, 8 mol% CO2, and 28 mol% C3H8) with heavy oil. Then, a theoretical methodology was employed to estimate the phase behavior of the heavy oil–solvent mixture (HO–SM) systems with various mole fractions of SM. The experimental results show that as the mole fraction of SM increases, the saturation pressures and swelling factors of the HO–SM systems considerably increase, and the viscosities and densities of the HO–SM systems decrease. The heavy oil is upgraded in situ via asphaltene precipitation and SM dissolution. Therefore, the solvent-enriched oil phase at the top layer of reservoirs can easily be produced from the reservoir. The aforementioned results indicate that the SM has promising application potential for enhanced heavy oil recovery via solvent-based processes. The theoretical methodology can accurately predict the saturation pressures, swelling factors, and densities of HO–SM systems with various mole fractions of SM, with average error percentages of 1.77% for saturation pressures, 0.07% for swelling factors, and 0.07% for densities.

Keywords Heavy oil · Solvent mixture · PR-EOS · Phase behavior · Pressure–volume–temperature (PVT)

1 Introduction

The observed performance from many heavy oil reservoirs in Canada, China, and Venezuela, such as the Orinoco Oil Belt, Lindbergh, and Tuha Fields, has been significantly better than that from conventional heavy oil reservoirs during primary production processes (Abusahmin et al. 2017; Guan et al. 2008; Maini 1999; Sun et al. 2017a, 2019a, b; Zhou et al. 2016). One of the possible reasons for this difference is the in situ foamy oil occurrence (the solution gas is trapped in the oil phase as dispersed bubbles) (Du et al. 2016; Liu et al. 2017; Sun et al. 2014, 2018a; Zhang et al. 2015; Zhou et al. 2017). However, the gas bubbles eventually become a continuous free gas due to low reservoir pressure. Therefore, the foamy oil phenomenon disappears after primary production, resulting in low ultimate oil recovery factors (5%–15%) in heavy oil reservoirs (Shokri and Babadagli 2017; Soh et al. 2018).

Extensive studies have been performed to investigate the potential of solvent-based processes for heavy oil recovery after primary production. The recently examined solvent-based processes include vapor extraction (Kordestany et al. 2019; Talbi et al. 2008; Upreti et al. 2007), cyclic solvent injection (CSI) (Du et al. 2017; Sun et al. 2015; You et al. 2019), and solvent flooding (Luo et al. 2012, Zhang et al. 2019a). Generally, solvents (e.g., CO2, CH4, and C3H8) are used in a pure or mixed form during the aforementioned processes. The beneficial mechanisms of solvent-based processes mainly include enhanced oil swelling and oil mobility improvement, as well as reductions in interfacial tension (IFT), oil viscosity and density (Dong et al. 2013; Haddadnia et al. 2018a, b; Liu 2019; Yu and Shen 2008; Zhang et al. 2019b). All of these mechanisms can be determined by phase behavior data (Li et al. 2011, 2012a, 2013). Thus, obtaining accurate phase behavior data is vital in analyses of solvent-based processes.

The phase behavior of various solvents with heavy oil has been systematically investigated in previous
studies. Frauenfeld et al. (2002) measured the viscosities of CH$_4$, CO$_2$, and C$_3$H$_8$ in Lloydminster heavy oil (LHO) at 292.15 K and several pressures. Luo et al. (2007a) measured the viscosities and swelling factors of different LHO–C$_3$H$_8$ systems at pressures from 200 to 800 kPa. Badamchi-Zadeh et al. (2009) investigated the saturation pressures, liquid phase densities and viscosities of C$_3$H$_8$–Athabasca bitumen systems for temperatures from 283.15 to 323.15 K. Li et al. (2011) measured the swelling factors, saturation pressures, and viscosities of heavy oil–three solvent (C$_2$H$_6$, CO$_2$–C$_3$H$_8$, and CO$_2$–n-C$_4$H$_{10}$) systems at temperatures from 280.45 to 391.55 K.

A comprehensive experimental study of solvent and bitumen/heavy oil systems can be found in recent publications by SHARP research group at University of Calgary (Azinfar et al. 2018a, b; Haddadnia et al. 2018a, b; Nourozieh et al. 2014; Zirrahi et al. 2014, 2017a, b). Zirrahi et al. (2014) measured the solubility of CH$_4$, C$_2$H$_6$, and CO$_2$ in bitumen. Nourozieh et al. (2014) measured the densities of bitumen–pentane mixtures at temperatures varying from ambient up to 200 °C and at pressures up to 10 MPa. Zirrahi et al. (2017a, b) and Azinfar et al. (2018a, b) measured the phase behavior data including solubilities, densities, and viscosities of propane–bitumen (heavy oil) systems in wide ranges of temperature and pressure (including four temperatures of 50, 100, 150, and 186 °C, and pressures up to 8.1 MPa). Haddadnia et al. (2018a) and Azinfar et al. (2018a, b) reported the viscosity and solubility data of butane–Athabasca bitumen (up to 260 °C). In addition, they also investigated the thermos-physical properties of n-pentane–bitumen systems and n-hexane–bitumen systems at different temperatures (30–190 °C), pressures (2–8 MPa), and solvent mass fractions (0.05–0.5) (Haddadnia et al. 2018b).

Based on the aforementioned phase behavior data, many models were developed to predict thermos-physical properties of various solvents (such as CH$_4$, C$_2$H$_6$, CO$_2$, C$_3$H$_8$, and so on) in bitumen (heavy oil). The PR cubic equation of states (EOS) and Krichevsky–Ilinskaya equation were used in those models (Haddadnia et al. 2018b; Zirrahi et al. 2014). Recently, there are some studies that have developed reliable correlations or models for predicting thermos-physical properties of solvent–heavy oil (bitumen) systems by artificial intelligence. For example, Ahmadi et al. (2018) developed a predictive model based on the least-squares support vector machine (LS-SVM) to calculate the CO$_2$ and heavy oil swelling factor and found that the LS-SVM is a straightforward and accurate method to determine the oil swelling factor with negligible uncertainty. Al-Gawfi et al. (2019) developed three generalized correlations using symbolic regression based on genetic programming to predict the solubility, density, and viscosity of light hydrocarbon (C$_1$ to C$_3$) and bitumen mixtures.

There are a limited number of studies on the phase behavior of heavy oil–solvent mixture (HO–SM) systems. Li et al. (2013) experimentally determined the swelling factors of LHO–C$_3$H$_8$–n–C$_4$H$_{10}$ systems with different solvent mole fractions at high temperatures and pressures. Yang et al. (2014) determined the saturation pressures and swelling factors of LHO–CH$_4$–C$_3$H$_8$ systems using the PVT tests and a PR-EOS model, respectively. Li et al. (2013) proposed a generalized methodology to predict the swelling factors of LHO–CO$_2$–C$_3$H$_8$ and LHO–CO$_2$–n–C$_4$H$_{10}$ systems under reservoir conditions. Yamchi et al. (2018) added propane to the mixtures of butane–bitumen systems and determined the changes in thermos-physical properties. The results showed that the addition of propane resulted in lighter extracted phase and an increase in the saturation pressure of solution. Zhao et al. (2019) measured the saturation pressures of three HO–SM (CH$_4$CO$_2$, CH$_4$–C$_3$H$_8$, and CO$_2$–C$_3$H$_8$) systems by constant composition expansion tests.

According to the aforementioned studies in the literature, large quantities of phase behavior data are available for various solvents and LHO/bitumen (Athabasca and Cold Lake). However, many heavy oil reservoirs, such as those in the Orinoco Oil Belt, Venezuela, have higher reservoir pressures (8650 kPa) than that in Canada (1800–2400 kPa). The solvents or solvent mixtures for LHO and Athabasca and Cold Lake bitumen investigated in the literature are not suitable for those heavy oil reservoirs with high pressures. Therefore, it is crucial for finding an effective injection solvent mixture for solvent-based processes that is applicable to heavy oils with higher pressures.

Compared with commonly used CO$_2$ and CH$_4$, C$_3$H$_8$ has higher solubility in heavy oil. Therefore, the addition of C$_3$H$_8$ increases the contribution of the solvent mixture (SM) to increased oil swelling and reductions in oil viscosity and IFT (Ahadi and Torabi 2018). However, due to the relatively low dew point pressure of C$_3$H$_8$ (C$_3$H$_8$ is easily liquefied in the high-pressure reservoir conditions, resulting in larger solvent consumption (Fig. 1).

Therefore, the addition of produced gas (generally, the main components are CO$_2$ and CH$_4$) as a carrier gas allows the SM to be gas phase under reservoir conditions, thereby avoiding the C$_3$H$_8$ liquefaction and reducing the C$_3$H$_8$ consumption. In addition, produced gas can be obtained from the primary production of heavy oil reservoirs, and its use improves the availability and economic efficiency of solvent-based processes. Therefore, the SM (CH$_4$, CO$_2$, and C$_3$H$_8$) can fully utilize the advantages of C$_3$H$_8$ and produced gas, which have significant potential in heavy oil reservoirs after primary production.

To the best of our knowledge, the phase behavior of HO–SM systems (CH$_4$, CO$_2$, and C$_3$H$_8$) has not yet been investigated although it is crucial for revealing the mechanisms of the SM and designing effective solvent-based
processes. It is still lack of basic experimental data for phase behavior of HO and SM (CH4, CO2, and C3H8) systems. In addition, the theoretical methodology is still not available to estimate the phase behavior of HO–SM systems. In this paper, a novel experimental procedure is proposed to investigate the phase behavior of HO–SM systems. First, the saturation pressures, swelling factors, viscosities, and densities of the systems with various mole fractions of SM ranging from 10.29 to 33.54 mol% were measured using PVT tests. Then, quantitative analyses were performed to characterize the HO–SM systems at different layers in the PVT cell by measuring the viscosities and densities at different saturation pressures. Finally, a theoretical methodology was proposed to predict the theoretical phase behavior of the HO–SM systems with various mole fractions of SM.

2 Experimental section

2.1 Experimental materials

A heavy oil (HO) sample from the Orinoco Oil Belt, Venezuela (VHO), was used in this study. The physical properties of the VHO sample are presented in Tables 1 and 2, and the compositional analysis results obtained by the Schlumberger Venezuela are shown in Table 3.

The densities of the HO sample were obtained at experimental pressures from 100 to 10,000 kPa and experimental temperatures from 298.15 to 338.15 K with a density meter, and the viscosities of the HO sample at temperatures from 298.15 to 368.15 K and atmospheric

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**Table 1** SARA analysis of the VHO sample

| Saturates, wt% | Aromatics, wt% | Resins, wt% | Asphaltenes, wt% |
|---------------|----------------|-------------|------------------|
| 19.8          | 51.2           | 18.9        | 8.8              |

**Table 2** VHO sample properties

| Physical properties | Values |
|---------------------|--------|
| Initial gas–oil ratio, m³/m³ | 15     |
| Formation volume factor at saturation pressure, m³/m³ | 1.173  |
| Bubble point pressure, MPa | 4.95   |
| Acid number, mg/g | 4.95   |
| Base number, mg/g | 13.57  |

**Table 3** Compositional analysis results of the VHO sample (in mol%)

| Components | Reservoir oil | Produced gas | Produced oil |
|------------|--------------|--------------|--------------|
| N2         | 0.13         | 0.52         | 0.00         |
| CO2        | 2.80         | 10.82        | 0.00         |
| C1         | 22.43        | 86.72        | 0.00         |
| C2         | 0.08         | 0.31         | 0.00         |
| C3         | 0.04         | 0.17         | 0.00         |
| C4         | 0.04         | 0.15         | 0.00         |
| C5         | 0.12         | 0.13         | 0.00         |
| C6         | 0.49         | 1.05         | 0.29         |
| C7         | 0.25         | 0.09         | 0.30         |
| C8         | 0.32         | 0.03         | 0.42         |
| C9         | 0.50         | 0.01         | 0.67         |
| C10        | 0.90         | 0.00         | 1.21         |
| C11        | 1.53         | 0.00         | 2.06         |
| C12        | 1.84         | 0.00         | 2.48         |
| C13        | 1.96         | 0.00         | 2.65         |
| C14        | 2.13         | 0.00         | 2.87         |
| C15        | 2.24         | 0.00         | 3.02         |
| C16        | 2.38         | 0.00         | 3.21         |
| C17        | 2.39         | 0.00         | 3.22         |
| C18        | 2.50         | 0.00         | 3.37         |
| C19        | 2.56         | 0.00         | 3.45         |
| C20        | 2.51         | 0.00         | 3.38         |
| C21        | 2.31         | 0.00         | 3.11         |
| C22        | 2.18         | 0.00         | 2.94         |
| C23        | 2.08         | 0.00         | 2.81         |
| C24        | 1.91         | 0.00         | 2.58         |
| C25        | 1.85         | 0.00         | 2.49         |
| C26        | 1.65         | 0.00         | 2.23         |
| C27        | 1.50         | 0.00         | 2.03         |
| C28        | 1.38         | 0.00         | 1.86         |
| C29        | 1.25         | 0.00         | 1.69         |
| C30+       | 33.76        | 0.00         | 45.54        |

Molecular weight 418.76 20.19 557.81

Molecular weight of C30+ 916.89 0.00 916.89
pressure were measured using a rolling-ball viscometer (see Experimental section for detailed information about the viscometer and density meter). The measured viscosities and densities of the HO sample are shown in Fig. 2.

Figure 2 also shows that the viscosities and densities calculated with Eqs. (1)–(5) are in good agreement with the measured results (Li et al. 2011; Moriyoshi et al. 1977).

\[
\log_{10} \left[ \log_{10}(\mu) \right] = -3.2943 \log_{10}(T) + 8.9092
\]

where \( T \) is the experimental temperature in K; \( \rho \) is the HO density in kg/m\(^3\); \( P \) is the experimental pressure in kPa; and \( \mu \) is the HO viscosity in mPa s.

The parameters \( \rho_0, \beta, \) and \( B \) are related to temperature based on the following equations (Moriyoshi et al. 1977):

\[
\rho_0 = a_1 + a_2 T + a_3 T^2
\]

\[
B = b_1 + b_2 T^{-1} + b_3 T^{-2}
\]

\[
\beta = c_1 + c_2 T
\]

The parameter values \( a_1, a_2, a_3, b_1, b_2, b_3, c_1 \) and \( c_2 \) shown in Eqs. (3)–(5) are \( 9.414 \times 10^2, 9.393 \times 10^{-1}, -2.500 \times 10^{-3}, 4.101, -1.752 \times 10^3, 2.636 \times 10^5, -3.271 \times 10^{-3}, \) and \( 3.545 \times 10^{-5} \), respectively, which were determined based on a regression analysis of the density data.

The gases CO\(_2\), CH\(_4\), and C\(_3\)H\(_8\) used in this study had a purity of 99.99%, and the solvent mixture (SM) was generated by mixing CH\(_4\), CO\(_2\), and C\(_3\)H\(_8\) with a molar fraction ratio of 64:8:28.

### 2.2 Experimental setup

A schematic diagram of the experimental setup is shown in Fig. 3.

The main component of the experimental setup is a PVT system (FY-PVT-1, FEIYU petroleum technology Co. Ltd., China). The PVT cell in a thermal jacket has a length of 35.00 cm and an inner diameter of 3.57 cm. The PVT cell was made of stainless steel (1Cr18Ni9Ti), and it can withstand up to 70,000 kPa of pressure and temperatures up to 453.15 K. The volume and pressure of the PVT cell are adjusted with a high-pressure pump that is controlled by a control panel attached to a support frame. A high-precision test gauge (Druck Ltd., UK) with an accuracy of ±1 kPa is applied to measure the pressure in the PVT cell, and a temperature controller (SRS-13, Shimadzu, Japan) with an accuracy of ±0.1 K is used to control the temperature in the PVT cell.

Compared with light or medium oils, PVT tests of heavy oils are more time-consuming because a longer time is needed for gas liberation or dissolution in heavy oil. A special rotary mixing system is designed to accelerate the PVT tests for the HO–SM systems. The rotary mixing system mainly consists of an agitator fully immersed in the fluid and an electromotor located at the bottom of the support frame (Fig. 3). Considering the high viscosity of the HO, the agitator is made of high-density wolfram (15.63 g/cm\(^3\)). The electromotor renders the central axis rotates, and then the PVT cell rotates up and down around the central axis in the vertical plane. Due to the movement of the PVT cell, the agitator moves up and down in the PVT cell during this process, which efficiently mixes the heavy oil and the SM at the selected pressures and temperatures. The rotation speed
The heavy oil and single solvents were first stored in sample cylinders and then used to form the SM in the transfer cylinder with a thermal jacket. A high-precision pump was employed to transfer the SM from the cylinder to the PVT cell.

A rolling-ball viscometer (CHY-V, SHITIAN electronics Co. Ltd., China) was applied to measure the viscosities of the HO–SM systems. The temperature of the viscometer was maintained via a thermal jacket with an accuracy of ±0.1 K. The measurement limits of the viscometer were 70,000 kPa and 473.15 K, respectively. The viscometer was calibrated using standard viscosity liquids (National Institute of Metrology, China) in the range from 2368.2 to 9308 mPa s with a precision of ±5 mPa s. A density meter (Anton Paar, Austria) was used to determine the densities of the HO–SM systems. This meter was calibrated using water at pressures between 690 and 6900 kPa and temperatures between 283.15 and 323.15 K, and the precision was ±0.001 g/cm³.

2.3 Experimental procedures

Before each test, the PVT cell was cleaned with kerosene and evacuated using a vacuum pump. Subsequently, a desired volume of HO was transferred into the PVT cell. The HO mass was determined from Eq. (2). According to the mole fraction of the SM (10.29 mol%) in the studied system and the composition of the SM (64 mol% CH₄, 8 mol% CO₂, and 28 mol% C₃H₈), the amounts of CH₄, CO₂, and C₃H₈ were calculated based on the equation of state for real gases (Li et al. 2012b). Then, the required amounts of CH₄, CO₂, and C₃H₈ were introduced into the transfer cylinder to generate the SM. The temperature of the transfer cylinder was kept at the reservoir temperature (327.35 K) for 24 h by the thermal jacket to confirm that the PVT cell and transfer cylinder reached the same temperature. Finally, the SM was injected into the PVT cell to constitute the HO–SM system with the desired mole fraction of SM using a high-precision pump.

To accelerate the equilibrium process between the HO and SM, the PVT cell pressure was increased to 15,000 kPa, which was more than the reservoir pressure (8650 kPa). Because the continuous depressurization method is reported to be less accurate than the discrete depressurization method (Badamchi-Zadeh et al. 2009; Li et al. 2011), the discrete depressurization method was applied to determine the saturation pressure of the HO–SM system in this study, even if it took a longer time to maintain the system in equilibrium. In this process, the PVT cell was incrementally depressurized by increasing the PVT cell volume. At the end of each depressurization step, the rotary mixing system was set to sufficiently stir the HO and SM until the PVT cell pressure stayed constant over a period of time. The equilibrium between the HO and SM

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**Fig. 3** Schematic diagram of experimental tests
was considered to be reached. Consequently, the saturation pressure and the corresponding system volume were determined from the transition point on the pressure–volume \((P-V)\) curve. Finally, the swelling factor of the HO–SM system with a desired mole fraction of SM was calculated as the volume of HO–SM system at the saturation pressure divided by the HO volume at atmospheric pressure. The repeatabilities of the saturation pressure and volume measurements were found to be ±2% and ±3%, respectively.

After the saturation pressure and swelling factor measurements, the system in the PVT cell was compressed and re-equilibrated at the saturation pressure. Then, a certain amount of the mixture at the top of the PVT cell was passed to the density meter and viscometer, and the viscosity and density of the system at the saturation pressure were determined.

In addition, to examine the change of viscosities and densities of the mixture at different depths after the in situ heavy oil upgrading process, the remaining mixture (150 m³) in the PVT cell was kept at the saturation pressure without stirring for 45 days. Then, the mixture was artificially divided into three layers (around 50 m³) shown in Fig. 3, according to the volume requirements of the falling ball viscometer and density meter. Overall, there is no clearly defined boundary between any two layers, and such a division is subjective. In this process, the oil of the top layer was divided into three oil samples (the volume of each sample was around 15 m³), and they were introduced into the falling ball viscometer and the density meter for measuring their viscosity and density at saturation pressure in turn. The average values of the three measurements represent the viscosity and density of the top layer. Following the abovementioned process, the viscosities and densities of middle layer and bottom layer at deeper locations of the PVT cell were measured one by one. The repeatabilities of the density and viscosity measurements were found to be ±2% and ±3%, respectively.

In this study, HO–SM systems with varied solvent mole fractions (10.29–33.54 mol%) were prepared and tested at 327.35 K following the aforementioned procedure. After the measurements of the saturation pressures, swelling factors, viscosities, and densities at the lowest mole fraction of SM were obtained, the mole fraction of SM was increased, and the measurements were repeated at the higher mole fraction of SM. Thus, the effects of the mole fraction of SM on the phase behaviors of the HO–SM systems were systematically investigated.

### 3 Theoretical section

#### 3.1 Equation of state

The PR-EOS model is used and given by Eqs. (6)–(8) (Peng and Robinson 1976):

\[
P = \frac{TR}{V - b} - \frac{a}{b(V - b) + V(V + b)} \tag{6}
\]

\[
a = 0.457235 \frac{R^2T^2}{PC_c} \tag{7}
\]

\[
b = \frac{0.0777969RT_c}{PC_c} \tag{8}
\]

where \(R\) is the universal gas constant in kPa m³/(K kmol); \(a\) is the attraction parameter; \(V\) is the molar volume of the HO–SM system in m³/kmol; \(b\) is a constant related to critical properties in m³/kmol; \(P_c\) is the critical pressure of each component in the HO–SM system in kPa; \(\alpha\) is the dimensionless function; and \(T_c\) is the critical temperature of each component in the HO–SM system in K.

The commonly used \(\alpha\) function is a Soave-type function (Peng and Robinson 1976):

\[
\alpha = \left[1 + \left(0.37464 + 1.54226\omega - 0.26992\omega^2\right)(1 - T_r^{0.5})\right]^2 \tag{9}
\]

where \(\omega\) and \(T_r\) are the acentric factor and reduced temperature of each component in the HO–SM system, respectively.

However, according to previous studies, the following \(\alpha\) function can provide a more accurate estimation of the saturation pressure compared with the Soave-type function (Li and Yang 2011):

\[
a = \exp\left\{\left(0.13280 - 0.05052\omega + 0.25948\omega^2\right)(1 - T_r) + 0.81769 \ln \left[1 + (0.31355 + 1.86745\omega - 0.52604\omega^2)(1 - \sqrt{T_r})\right] \right\} \tag{10}
\]

The studied HO–SM system is a mixture. Therefore, the \(a\) and \(b\) used in Eq. (6) should be calculated by the following mixing rules.

\[
a = \sum_{i=1}^{n} \sum_{j=1}^{n} x_ix_j(1 - \delta_{ij}) \sqrt{a_ia_j} \quad (i, j = 1, 2, 3, 4) \tag{11}
\]

\[
b = \sum_{i=1}^{n} x_ib_i \tag{12}
\]

where \(n\) is the component number in the HO–SM system; \(\delta_{ij}\) is the empirically determined binary interaction parameter (BIP) between components \(i\) and \(j\); \(x\) is the mole fractions.
of each component in the HO–SM system. In this study, the heavy oil is characterized as a single pseudocomponent, HO as its utilization in PR-EOS model can accurately predict the phase behavior of solvent–heavy oil systems, as proven by previous research (Li and Yang 2013; Li et al. 2011, 2012c; Yang et al. 2014). Therefore, totally four components (C3H8, CO2, CH4, and HO) are used to characterize the HO–SM system in this study.

### 3.2 BIP correlations

The BIP between any two components must be calculated based on Eq. (11). Because C3H8, CO2, and CH4 are pure components, the BIPs for the C3H8-CO2 pair (δ12), C3H8-CH4 pair (δ13), and CO2-CH4 pair (δ23) can be determined by the following equation (Chueh and Prausnitz 1967; Oellrich et al. 1981):

\[
d_{ij} = 1 - \left[ 2 \sqrt{\frac{V_{c1}^{1/3} V_{c2}^{1/3}}{V_{c1}^{1/3} + V_{c2}^{1/3}}} \right]^{1.2} \quad (i, j = 1, 2, 3, 4) \tag{13}
\]

where \( V_c \) is the critical molar volume of each component in the HO–SM system.

The BIPs for the C3H8-HO pair (δ14), CO2-HO pair (δ24), and CH4-HO pair (δ34) are determined with the following correlations that were developed by fitting experimental solubility data, as recommended by several researchers (Li and Yang 2013; Li et al. 2011, 2012c; Yang et al. 2014).

- C3H8-HO pair: \( \delta_{14} = -0.4560 \frac{T}{T_c} + 0.1817 \) \tag{14}
- CO2-HO pair: \( \delta_{24} = -0.5462 \frac{T}{T_c} - 0.4596\gamma - 0.0238\omega + 0.7523 \) \tag{15}
- CH4-HO pair: \( \delta_{34} = -0.8060 \frac{T}{T_c} - 0.8550\gamma - 0.0809\omega + 1.1880 \) \tag{16}

where \( \gamma, \omega, \) and \( T_c \) are the specific gravity, acentric factor, and critical temperature of HO, respectively.

The critical properties (\( T_c, \omega, \) and \( P_c \)) of CO2, CH4, and C3H8 are available in the literature (Li and Yang 2013; Yang et al. 2014), and the critical properties of HO can be obtained from the equations in “Appendix”.

### 3.3 Swelling factor and density

The swelling factor is of great importance for evaluating the validity of a solvent during solvent-based processes. Generally, oil swelling caused by solvent dissolution can increase heavy oil saturation and mobility in porous media (Li et al. 2013). Therefore, a larger swelling factor is beneficial for the performance of a solvent-based process.

The swelling factor, \( f \), of the HO–SM system is defined as follows (Li et al. 2011):

\[
f = \frac{V_2}{V_1 (1 - S)} \tag{17}
\]

where \( S \) is the mole fraction of the SM, \( V_1 \) is the molar volume of HO at 101.32 kPa in m\(^3\)/kmol, and \( V_2 \) is the molar volume of the HO–SM system at the saturation pressure and reservoir temperature in m\(^3\)/kmol.

In this study, three volume translation methods were used to improve swelling factor prediction. Notably, these methods have no impact on the prediction of saturation pressures (Peneloux et al. 1982). The corrected value of \( V_2 \) is calculated as follows:

\[
V_2 = V - \sum_{i=1}^{n} x_i C_i \tag{18}
\]

\[
C = s \times b \tag{19}
\]

where \( C \) is the correction parameter; and \( s \) is the shift parameter. This shift parameter can be predicted by the equation developed by Jhaveri and Youngren (1988):

\[
s = 1 - \frac{d}{M e} \tag{20}
\]

where \( M \) is the molecular weight and \( e \) and \( d \) are constants equal to 0.1823 and 2.258, respectively.

Furthermore, Peneloux et al. (1982) and Twu and Chan (2009) proposed two correlations to estimate \( C \), as shown in Eqs. (21) and (22):

\[
C = 0.40768 \left( \frac{RT_c}{P_c} \right) (0.29441 - Z_{RA}) \tag{21}
\]

\[
C = 0.406501 \left( \frac{RT_c}{P_c} \right) (0.260484 - Z_{RA}) \tag{22}
\]

where \( Z_{RA} \) is the Rackett’s compressibility factor of VHO, which can be expressed as follows:

\[
Z_{RA} = \left( \frac{MP_c}{RT_c} \right)^{1/\left[1+(1-T_c)/P_c\right]^2} \tag{23}
\]

The aforementioned volume translation methods were evaluated based on their accuracy for predicting the swelling factors of the HO–SM systems with various mole fractions of SM.
The density of the HO–SM system can be calculated as follows:

\[
\rho_{\text{mix}} = \frac{n_{o}M_{o} + n_{g}M_{g}}{V_{o}}
\]  

(24)

where \(\rho_{\text{mix}}\) is the density of HO–SM system in kg/m\(^3\); \(n_{o}\) and \(n_{g}\) are the amount of substance of HO and SM in kmol; \(M_{o}\) and \(M_{g}\) are the molecular weight of HO and SM in kg/kmol; \(V_{o}\) is the volume of HO in m\(^3\).

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**Fig. 4** Measured \(P–V\) curves of the HO–SM systems with varied mole fractions: a 10.29 mol% SM, b 18.66 mol% SM, c 26.59 mol% SM, d 31.50 mol% SM, and e 33.54 mol% SM.
4 Results and discussion

4.1 Discussion on the experimental findings

Figure 4 demonstrates the $P$–$V$ curves of the HO–SM systems with varied mole fractions of SM.

The saturation pressures and the corresponding volumes were obtained based on the transition points of the curves shown in Fig. 4. The measured saturation pressures and calculated swelling factors at different mole fractions of SM are plotted in Fig. 5.

As shown in Fig. 5, the saturation pressures considerably increase as the mole fraction of SM increases, implying that a high pressure is needed if a large amount of injected SM is dissolved in the HO during solvent-based processes. In addition, the high values of the measured saturation pressures indicate that the dissolved SM is released from the oil phase during the production period of CSI processes, providing sufficient energy for driving solution gas. Moreover, because a high mole fraction of SM exists even at low pressures (Fig. 5), the viscosity and density of the system remain low during the late stage of the production period in CSI processes.

Figure 5 also shows that the measured swelling factor curve increases almost linearly with increasing mole fraction of SM, which suggests that the dissolution of the SM leads to the swelling of the HO. The HO swells 1.034 times due to the dissolution of the 33.54 mol% SM. However, the magnitude of swelling is not as drastic as that for light oils, which can swell by approximately 10% of their original volume due to significant solvent dissolution (Luo et al. 2007a). A large swelling factor is beneficial for the performance of solvent-based processes because the enhanced oil swelling effect can lead to an increased relative permeability in the oil phase and enhanced mobile oil saturation in heavy reservoirs.

Viscosity and density reduction is an important mechanism of solvent-based processes for enhancing heavy oil recovery. Solvent can significantly reduce the oil viscosity and density, even at a relatively low pressure (Kariznovi et al. 2011). The measured viscosities and densities of the HO–SM systems at saturation pressures are shown in Fig. 6. Figure 6 clearly shows that the viscosities and densities of the systems generally decrease as the mole fraction of SM in each system increases, as expected. For example, the viscosity and density of the HO–SM system were reduced from 16,105.40 mPa s and 0.9812 g/cm$^3$ to 5585.70 mPa s and 0.9706 g/cm$^3$ with the dissolution of a 33.54% SM. Therefore, it can be inferred that increasing the dissolution of the SM at a high pressure or a large gas injection volume is beneficial for performance in solvent-based processes.

Figure 7 presents the viscosities and densities of the HO–SM systems in the three different layers at saturation pressures after 45 days of equilibrium.

As shown in Fig. 7, regardless of the mole fraction of the SM, the viscosities and densities of the HO–SM systems increase with increasing depth in the PVT cell. For example, when the mole fraction of the SM is 10.29%, the viscosity and density of the bottom layer are 11,174.70 mPa s and 0.9841 g/cm$^3$, respectively, which are approximately 1.25 and 1.01 times those values of 8951.06 mPa s and 0.9707 g/cm$^3$ of the top layer. Therefore, the HO–SM systems become heterogeneous and exhibit different physicochemical properties inside the PVT cell after 45 days of equilibrium.

This finding can contribute the in situ heavy oil upgrading under the experimental conditions. After the HO–SM systems were in a static state for 45 days, some light components were continuously extracted from the HO at the top of the PVT cell. Moreover, the SM concentration in the HO resulted in asphaltene precipitation, and asphaltene was deposited at the bottom of the PVT cell. Therefore, the top layer becomes an oil phase enriched with solvent.
Additionally, the middle layer is HO with dissolved SM, and the bottom layer becomes HO with heavy components.

Figure 7 shows that the viscosities and densities of the top layer gradually decrease as the mole fraction of the SM increases. However, the viscosities and densities of the middle and bottom layers initially decrease, and then a slight increase is observed. This result occurs because the top layer is mainly light components. When the mole fraction of the SM increases, more dissolution of the SM in the solvent-enriched oil phase causes decreases in the viscosity and density of the top layer. However, for the mixtures in the middle and bottom layers, the precipitation of asphaltene significantly increases when the mole fraction of SM exceeds 31.50 mol%, which contributes to slight increases in the viscosity and density of HO in the bottom layer, as shown in Fig. 7. These observations are consistent with the results obtained in previous studies involving asphaltene precipitation tests (Luo et al. 2007b).

In a field application, if a large amount of SM is injected into the HO reservoir, the generated asphaltene precipitation can result in the in situ heavy oil upgrading. When an additional SM is injected and dissolved into the upgraded HO, the viscosity and density of the in situ upgraded HO can be further reduced. Therefore, the HO in the top layer can be easily produced via solvent-based processes. However, the HO in the bottom layer, which mainly contains heavy components, including asphaltene deposits, is left behind in the HO reservoir.

Overall, the phase behavior data for the HO–SM systems demonstrate that the SM exhibit promising application potential for heavy oil recovery in solvent-based processes. The dominant mechanisms include combined viscosity and density reduction, solution gas driving, oil swelling, and the in situ upgrading of heavy oil.

### 4.2 Discussion on the prediction results

The calculated critical properties for all the components and BIPs for all the component pairs are shown in Tables 4 and 5. The saturation pressures, swelling factors, and densities

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Table 4

| Component      | \(T_c\)  | \(P_c\)  | \(Z_{RA}\) | \(\omega\) | \(V_c\)  | References    |
|----------------|--------|--------|----------|--------|--------|---------------|
| \(\text{C}_3\text{H}_8\) | 369.80 | 4246   | 0.2763   | 0.1520 | 0.2030 | Li et al. (2011) |
| \(\text{CO}_2\)     | 304.14 | 7378   | 0.2736   | 0.2238 | 0.0940 | Li et al. (2011) |
| \(\text{CH}_4\)     | 190.55 | 4600   | 0.2876   | 0.0080 | 0.0990 | Yang et al. (2014) |
| Heavy oil (VHO)     | 951.53 | 1010   | 0.2301   | 1.1556 | N/A    | "Appendix"    |

Table 5

| Component       | \(\text{C}_3\text{H}_8\) | \(\text{CO}_2\) | \(\text{CH}_4\) | Heavy oil (VHO) |
|-----------------|-----------------|----------------|---------------|-----------------|
| \(\text{C}_3\text{H}_8\) | 0               | 0.1350         | 0.0085        | 0.0248          |
| \(\text{CO}_2\)     | 0.1350          | 0              | 4.5 \times 10^{-5} | 0.0859          |
| \(\text{CH}_4\)     | 0.0085          | 4.5 \times 10^{-5} | 0             | 0.0217          |
| Heavy oil (VHO)    | 0.0248          | 0.0859         | 0.0217        | 0               |
for the HO–SM systems with varied mole fractions of SM were predicted by the PR-EOS model with Eqs. (6), (17), and (24) (Figs. 5 and 6). As shown in Figs. 5 and 6, the measured and predicted saturation pressures, swelling factors, and densities exhibit good agreement.

To quantitatively evaluate the accuracy of the predicted saturation pressures, swelling factors, and densities, three statistical methods were utilized: linear regression analysis, average error percentage (\(AEP = \frac{1}{n} \sum_{i=1}^{n} \left| \frac{X_{i}^{\text{pred}} - X_{i}^{\text{meas}}}{X_{i}^{\text{meas}}} \right|\), where \(n\) is the number of measured mole fractions of SM and \(X_{i}^{\text{pred}}\) and \(X_{i}^{\text{meas}}\) are the predicted and measured saturation pressure, swelling factor, or density values, respectively), and analysis of variance tables (difference between measured and predicted values), and analysis of variance tables (difference between measured and predicted values).

As shown in Fig. 8 and Table 6, the slope (\(A\)) of the linear correlation \((y = Ax + B)\) between predicted and measured saturation pressures is close to 1, and the correlation coefficient is 1.00. In addition, the \(AEP\) value between the measured and predicted results is 1.77%. All the results clearly indicate that good agreement exists between the measured and predicted results.

The calculated analysis of variance table for saturation pressures is shown in Table 7, where \(DF\) is the degrees of freedom, \(SS\) is the sum of the squares, \(MS\) is the mean squares, \(F\) is a variance-related statistical parameter. A detailed description of the parameters in Table 7 can be found elsewhere (Sun et al. 2017b, 2018b; Zendehboudi et al. 2011). As shown in Table 7, the calculated \(F\) value is 64,619.31, which is much higher than the critical value of \(F_{0.05}\) (7.71) and provides strong evidence against the null hypothesis. Therefore, all statistical analyses confirmed the suitability of the prediction method used in this study.

To evaluate the effects of volume translation methods on the prediction of swelling factors, the swelling factors were calculated using the Jhaveri method (Eq. 19), Peneloux method (Eq. 21), Twu method (Eq. 22), and Twu-Jhaveri method, respectively. The Twu-Jhaveri method is a combined calculation method in which the correction factors for the solvent components (\(\text{CH}_4, \text{CO}_2, \text{C}_3\text{H}_8\)) and HO component are calculated by the Jhaveri method and Twu method, respectively.

The predicted swelling factors are shown in Fig. 9. Notably, the Twu-Jhaveri method is the most reliable method for predicting the swelling factors of the HO–SM systems in comparison with the other three methods. The values

![Fig. 8](#) Regression for measured and predicted saturation pressures

![Fig. 9](#) Regression results for measured and predicted swelling factors

| Table 6 Regression for the measured and predicted saturation pressures |
|-------------------------|---------|---------|---------|-------------------|-------------------|------------------|
| Coefficient | Value | Upper 95% | Lower 95% | Standard error | Correlation coefficient | AEP, % |
| \(A\) | 1.03 | 1.04 | 1.01 | 0.00 | 1.00 | 1.77 |
| \(B\) | -0.10 | -0.03 | -0.16 | 0.02 |

| Table 7 Analysis of variance table for regression analysis of the measured and predicted saturation pressures |
|-------------------------------------|-----|-----|-----|-----|
| Source | DF | SS  | MS  | \(F\) |
| Regression | 1  | 27.78 | 5.56 | 64,619.31 |
| Residual | 3  | 0.00  | 0.00 |  |
| Total | 4  | 27.78 |  |

![Table 7](#) Analysis of variance table for regression analysis of the measured and predicted saturation pressures
predicted by the other methods are lower than the measured values, especially those obtained by the Jhaveri method and the Peneloux method.

The statistical analysis results show that the slope (A) of the linear correlation obtained by the Twu-Jhaveri method is 0.97 (Table 8 and Fig. 9). The slopes of the linear correlations obtained by the Jhaveri method, Peneloux method, and Twu method are 1.95, 0.68, and 0.85, respectively. A comparison of the slope values indicates that the value obtained by the Twu-Jhaveri method is 1.00, which implies that the Twu-Jhaveri method adequately reflects the experimental data. In addition, the linear correlations between the predicted and measured data obtained by the four methods all have high correlation coefficient values, and the F values are much greater than the corresponding $F_c$ value (7.71), which highlights the reliability of the regression results (Table 9).

In addition, Table 8 shows that the swelling factors can be reproduced by the Twu-Jhaveri method with an AEP value of 0.07%, and AEP values of 12.33%, 19.36%, and 0.32% for swelling factors were obtained by the Jhaveri method, Peneloux method, and Twu method, respectively.

As shown in Fig. 10 and Table 10, the statistical analysis results show that the slope (A) of the linear correlation between predicted and measured densities is 0.91, the correlation coefficient is 0.99, and the AEP value between the measured and predicted results is 0.07%. As shown in Table 11, the calculated $F$ value is 308.79, which is much higher than the critical value of $F_c$ (7.71), which highlights the reliability of the regression results.

Overall, the theoretical methodology based on the PR-EOS model, the $\alpha$ function, the selected BIP correlations, and the Twu-Jhaveri method can accurately predict the saturation pressures, swelling factors, and densities for the HO–SM systems with varied mole fractions of SM.
5 Conclusions

1. The theoretical methodology based on the PR-EOS model, the $\alpha$ function, the selected BIP correlations, and the Twu-Jhaveri method can accurately predict the saturation pressures, swelling factors, and densities of the SM (64 mol% $\text{CH}_4$, 8 mol% $\text{CO}_2$, and 28 mol% $\text{C}_3\text{H}_8$) for HO systems, with an average error percentages of 1.77% for saturation pressures, 0.07% for swelling factors, and 0.07% for densities.

2. The Twu-Jhaveri method was the most accurate method of predicting swelling factors for the HO–SM systems in comparison with the Jhaveri method, Peneloux method, and the Twu method. The swelling factors predicted by the Jhaveri method, Peneloux method, and the Twu method are lower than the experimental values, especially those obtained by the Jhaveri method and the Peneloux method.

3. As the mole fraction of SM increases, the saturation pressures and swelling factors considerably increase, and the viscosity and density of the HO–SM systems decrease. Therefore, an increase in the dissolution of the SM at high injection pressures or high gas injection volumes is beneficial for solvent-based processes.

4. The HO–SM systems become heterogeneous and exhibit different physicochemical properties after 45 days of equilibrium. The viscosity and density of the systems increase with increasing depth in the cell. The top layer is an oil phase enriched with solvent that can easily be recovered. However, the bottom layer mainly contains the heavy components, including asphaltene deposits, and is left behind in HO reservoirs.

5. As the mole fraction of the SM increases, the viscosity and density of the top layer gradually decrease. However, the viscosities and densities of the middle and bottom layers initially decrease, and then a slight increase is observed. Therefore, the in situ heavy oil upgrading significantly enhances the dissolution of SM in HO and reduces the viscosity and density of the HO.

6. The fundamental mechanisms associated with using the SM as a heavy oil recovery agent include driving the solution gas, reducing the viscosity and density, producing swelling effects, and upgrading the heavy oil in situ. The SM shows promising application potential for the enhancement of HO recovery in solvent-based processes.

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Appendix: Correlations of heavy oil properties

1. Critical temperature ($T_c$) (Kesler and Lee 1976; Li et al. 2013; Yang et al. 2014)

$$T_c = \frac{T_{cr}}{1.8}$$

2. Boiling point ($T_b$) = $\frac{T_{br}}{1.8}$

Table 10 Regression for the measured and predicted densities

| Coefficient | Value | Upper 95% | Lower 95% | Standard error | Correlation coefficient | AEP, % |
|-------------|-------|-----------|-----------|---------------|-------------------------|--------|
| $A$         | 0.91  | 1.03      | 0.72      | 0.05          | 0.99                    | 0.07   |
| $B$         | 91.59 | 275.86    | -32.68    | 48.47         |                         |        |

Table 11 Analysis of variance table for regression analysis of the measured and predicted densities

| Source   | DF  | SS   | MS   | F    |
|----------|-----|------|------|------|
| Regression | 1   | 37.39 | 7.48 | 308.79 |
| Residual  | 3   | 0.36 | 0.12 |      |
| Total     | 4   | 37.75|      |      |
\[ T_{CR} = 341.7 + 811\gamma + (0.4244 + 0.1174\gamma)T_{BR} + (0.4669 - 3.2623\gamma) \times 10^5T_{BR}^{-1} \]  
(27)

\[ T_{BR} = 1928.3 - (1.695 \times 10^3)M^{0.03522}\gamma^{-2.266} \times \exp \left[ - (4.922 \times 10^{-3})M - 4.7685\gamma + (3.462 \times 10^{-3})M\gamma \right] \]  
(28)

where \( T_b \) and \( T_c \) are the normal boiling point at 101.32 kPa and the critical temperature of HO in K.

2. Critical pressure \( (P_c) \) (Kesler and Lee 1976; Li et al. 2013; Yang et al. 2014)

\[ P_{CR} = \exp \left\{ \frac{8.3634 - 0.0566\gamma^{-1} - [(0.24244 + 2.2898\gamma^{-1} + 0.11857\gamma^{-2}) \times 10^{-3}]}{15.2518 - 15.6875T_{BR}^{-1} - 13.4721 \ln T_{BR} + 0.43577T_{BR}^6} \right\} \]  
(29)

\[ P_c = 6.895P_{CR} \]  
(30)

where \( P_{CR} \) is the critical pressure in psia.

3. Acentric factor \( (\omega) \) (Lee and Kesler 1975; Li et al. 2013; Yang et al. 2014)

If \( T_{BR} = T_{CR} \) \( T_{BR}^{-1} < 0.8 \)

\[ \omega = -\ln \left( \frac{P_c}{14.7} \right) - 5.92714 + 6.096487T_{BR}^{-1} + 1.28862 \ln T_{BR} - 0.169347T_{BR}^6 \]  
(31)

If \( T_{BR} = T_{CR} \) \( T_{BR}^{-1} > 0.8 \)

\[ \omega = -7.904 + 0.1352\frac{T_{BR}^{1/3}}{\gamma} - 0.007465 \left( \frac{T_{BR}^{1/3}}{\gamma} \right)^2 + 8.359T_{BR} \]  
\[ + \left( 1.408 - 0.01063\frac{T_{BR}^{1/3}}{\gamma} \right)T_{BR}^{-1} \]  
(32)

\[ T_b = T_{BR}/1.8 \]  
(33)

where \( T_b \) is the boiling point of HO at 101.32 kPa in K and \( T_{BR} \) is the reduced boiling point of HO.

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