CONSTRaining SOLAR FLare DIFFERENTIAL EMISSION MEASURES WITH EVE AND RHESSI

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ABSTRACT

Deriving a well-constrained differential emission measure (DEM) distribution for solar flares has historically been difficult, primarily because no single instrument is sensitive to the full range of coronal temperatures observed in flares, from $\lesssim 2$ to $\gtrsim 50$ MK. We present a new technique, combining extreme ultraviolet (EUV) spectra from the EUV Variability Experiment (EVE) onboard the Solar Dynamics Observatory with X-ray spectra from the Reuven Ramaty High Energy Solar Spectroscopic Imager (RHESSI), to derive, for the first time, a self-consistent, well-constrained DEM for jointly observed solar flares. EVE is sensitive to $\sim 2$–25 MK thermal plasma emission, and RHESSI to $\gtrsim 10$ MK; together, the two instruments cover the full range of flare coronal plasma temperatures. We have validated the new technique on artificial test data, and apply it to two X-class flares from solar cycle 24 to determine the flare DEM and its temporal evolution; the constraints on the thermal emission derived from the EVE data also constrain the low energy cutoff of the non-thermal electrons, a crucial parameter for flare energetics. The DEM analysis can also be used to predict the soft X-ray flux in the poorly observed $\sim 0.4$–5 nm range, with important applications for geospace science.

Key words: methods: data analysis – plasmas – radiation mechanisms: thermal – Sun: flares – Sun: UV radiation – Sun: X-rays, gamma rays

Online-only material: color figures

1. INTRODUCTION

Solar flares are powerful, explosive releases of magnetic energy, heating coronal plasma to tens of megaKelvin (MK) and accelerating electrons to hundreds of MeV. The physical mechanisms behind these processes are still poorly understood (see the review by Fletcher et al. 2011). While it is commonly accepted that much of the hot coronal plasma results from energy deposition in the chromosphere by non-thermal, accelerated particles—"chromospheric evaporation"—X-ray observations, e.g., by Yohkoh or the Reuven Ramaty High Energy Solar Spectroscopic Imager (RHESSI; Lin et al. 2002), suggest significant in situ coronal heating, as well (e.g., Masuda 1994; Masuda et al. 1998; Caspi & Lin 2010; Longcope & Guidoni 2011; Caspi et al. 2014), although the specific mechanism is still debated. Another crucial problem lies in identifying the low-energy cutoff/rollover of the accelerated electron spectrum, required by energetics (see the review by Holman et al. 2011) but typically poorly constrained as its observable manifestation—a rollover of the non-thermal bremsstrahlung spectrum (cf. Brown 1971)—is often obscured by thermal bremsstrahlung emission that dominates the photon spectrum up to $\sim 20$–35 keV in intense flares.

These questions remain, in large part, because of the difficulty in accurately and precisely characterizing the thermal electron population in flares. While flare analyses often employ the isothermal approximation (e.g., Lin et al. 1981; Garcia 1994; Feldman et al. 1996; Holman et al. 2003), decades of extreme ultraviolet (EUV) and X-ray observations with widely varying temperature sensitivities have shown that flare plasma exhibits a distribution of emission measure with temperature (the "differential emission measure," or DEM). Nonetheless, there have been only a few previous studies of solar flare DEMs (e.g., Dere & Cook 1979; McTiernan et al. 1999; Chifor et al. 2007; see also Trottet et al. 2011, who derived a four-component dis-
Figure 1. Left: DEMs derived from RHESSI (black) and EVE (red) data in isolation, with RHESSI dual-isothermal pre-fit (green plusses) and GOES XRS isothermal (blue plus) for comparison. The RHESSI DEM has limited resolution and poor constraints $\lesssim 10$ MK; the EVE DEM is better resolved but poorly constrained $\gtrsim 25$ MK. Center: observed RHESSI count rate spectrum vs. photon energy (black) and synthetic models derived from the EVE DEM (red) and GOES isothermal (blue); the high-temperature excess in the EVE-derived model overestimates observations by $\gtrsim 10\times$. The GOES-derived model agrees well at low energies, but indicates need for a higher-temperature component, consistent with the dual-isothermal fit. Right: arbitrarily truncating the EVE DEM at $\sim 30$ MK yields good model agreement with RHESSI at lower energies, consistent with the GOES-derived model, and suggests that instrument intercalibration is adequate.

(A color version of this figure is available in the online journal.)

2. METHOD DETAILS

The EVE suite comprises multiple instruments, including broadband photometers and high-resolution spectrometers that together measure solar EUV/SXR emission from $\sim 0.1$ to $\sim 105$ nm (Woods et al. 2012). We focus on the Multiple EUV Grating Spectrograph A (MEGS-A), a grazing-incidence dual-slit spectrograph measuring the spatially integrated solar spectral irradiance from $\sim 5$ to $\sim 37$ nm, with $\sim 0.1$ nm FWHM resolution and 10 s cadence (Hock et al. 2012). The MEGS-A wavelength range includes numerous spectral lines from various ion species with formation temperatures from $\sim 1$ to $\sim 25$ MK. The 9–15 nm range is particularly rich in high-temperature lines, including most of the strongest Fe XXII–XXIII lines (peak formation temperatures of $\sim 6$–15 MK; cf. Mazzotta et al. 1998): MEGS-A also observes Fe XXIV ($\sim 20$ MK) at 19.204 and 25.510 nm, and Fe XV ($\sim 2$ MK) and XVI ($\sim 3$ MK) at 28.416 and 33.541 nm, respectively, among many other lines. This wealth of temperature diagnostics enables the most precise determination of the flare DEM to date (Warren et al. 2013), especially from EVE observations, although they remain poorly constrained for temperatures $\gtrsim 25$ MK, where EVE has little or no temperature sensitivity (Figure 1). This can be somewhat mitigated by adding observations from the X-ray Sensor (XRS) on the Geostationary Operational Environmental Satellite (GOES), but the two-channel broadband XRS measurements—from which temperatures can also be inferred (e.g., White et al. 2005)—add minimal spectral information and extend the DEM validity only up to $\sim 30$ MK (Warren et al. 2013).

Significantly stronger constraints can be obtained using RHESSI (Lin et al. 2002), which measures solar hard X-rays (HXRs) below $\sim 0.4$ nm with $\lesssim 0.1$ nm FWHM resolution (quasi-constant $\sim 1$ keV when expressed in photon energy, variable $\sim \lambda^2/hc$ expressed in wavelength; Smith et al. 2002). RHESSI provides the most precise measurements of thermal X-ray continuum emission from plasmas with temperatures $\gtrsim 10$ MK, and is especially sensitive to the hottest part of the temperature distribution, up to $\gtrsim 50$ MK (e.g., Caspi et al. 2014). While RHESSI data can be used for DEM studies in isolation, RHESSI’s “temperature resolution” is broad because of the inherent “smoothing” of the bremsstrahlung emission mechanism; it is therefore difficult to recover a high-temperature-resolution DEM even from the high-temporal- and spectral-resolution continuum photon spectra, and the data often cannot sufficiently distinguish between a dual-isothermal model (delta-function DEM) or a continuous DEM (e.g., Caspi & Lin 2010; J. M. McTiernan & A. Caspi 2014, in preparation). The inferred DEM is also not well-constrained below $\sim 10$–$15$ MK, where the HXR yield is low and observations are typically limited by sensitivity or dynamic range (emission from the hotter parts of the DEM dominate the RHESSI spectra). EUV observations are required to provide these constraints.

For flares jointly observed by EVE and RHESSI, we have developed a technique to simultaneously analyze the full-Sun spectra from both instruments to derive a self-consistent DEM constrained over the entire range of coronal temperatures observed in flares, from $\sim 2$ MK to $\gtrsim 50$ MK. The lack of wavelength overlap and inherently different measurements make combining the two instruments’ data non-trivial; EVE data is best analyzed versus wavelength, and RHESSI versus photon energy. Because the mapping from physical parameters (e.g., temperature) to observables (e.g., irradiance) is not one-to-one—a photon of a given wavelength/energy could be emitted from many different temperatures or even from non-thermal electrons—particularly for resolution-blended lines in EVE or continuum emission in RHESSI, direct inversion of the combined data is ill-posed and prone to large uncertainties (cf. Craig & Brown 1976), especially in regions of poor signal-to-noise (e.g., between spectral lines in EVE, or at the longest wavelengths of RHESSI). Since the two instruments overlap in temperature sensitivity, however, we can overcome these issues through forward modeling—adopting a physical model, deriving observables, and optimizing the model parameters to maximize agreement—and we have developed a bridge to allow simultaneous analysis of both EVE and RHESSI data using the methods best suited to each instrument. Figure 1 shows that the instrument intercalibration—viz. the spectrum observed by one instrument vs. that predicted by the DEM derived independently from the other—is adequate.

Our technique combines and adapts the methods of Caspi & Lin (2010), Warren et al. (2013), and J. M. McTiernan & A. Caspi (2014, in preparation). A chosen flare is subdivided into analysis intervals of $\sim 60$–120 s, providing sufficient signal-to-noise while allowing study of dynamic processes. For each interval, the 10 s cadence EVE MEGS-A Level 2 (version 4) calibrated spectra, with default $\sim 0.02$ nm binning, are temporally averaged and the lines are isolated by subtracting the underlying continuum (whether solar or instrumental); the pre-flare...
Figure 2. Algorithm validation on artificial test input. Synthetic EVE (bottom) and RHESSI (top right) spectra are generated from the input DEM (top left), with added Poisson noise; the DEM is recovered by forward modeling. Orange highlights indicate the EVE wavelengths included in the fit; chromospheric lines are not synthesized. The recovered DEM (red), with constituent components shown (dotted), matches the input (black) well.

(A color version of this figure is available in the online journal.)

Our model DEM comprises multiple Gaussian components, equally spaced in log Te between 6.2 (~1.6 MK) and 7.8 (~63 MK), with fixed positions and widths but variable magnitudes; we have found that 10 Gaussians with FWHMs of ~0.19 in log Te sufficiently cover parameter space without overspecifying it. Each Gaussian is initialized to a random, uniformly distributed magnitude. Using the CHIANTI atomic database (v7.1; Dere et al. 1997; Landi et al. 2013), we compute high-resolution model spectra in the EUV (lines only) and X-ray (lines and continuum) ranges from the model DEM, then downsample to the respective instrument binning. To facilitate rapid calculations, the spectra are interpolated from a grid of CHIANTI emissivities precomputed as functions of temperature and wavelength/energy. The model EVE spectra are convolved with its nominal ~0.1 nm FWHM spectral resolution, while the model RHESSI spectra are convolved with its full instrument response and the pre-fit Fe/Fe–Ni and non-thermal components are added. We then use the Levenberg–Marquardt algorithm, as implemented in the IDL SolarSoft4 routine MPFIT (Markwardt 2009), to vary all the model parameters (Gaussian DEM magnitudes, Fe/Fe–Ni line fluxes, non-thermal power-law normalization, break energy, low-energy cutoff, and spectral indices) to iteratively minimize the deviations between the observed and model-computed spectra of both instruments simultaneously, as measured by the total reduced \( \chi^2 \) (the quadrature sum of the reduced \( \chi^2 \) for each instrument individually). For RHESSI, we use all statistically significant data from ~0.01 nm to ~0.25 nm. For EVE, we do not include the entire MEGS-A spectrum in the fit (see Figure 2), but rather focus on the line-rich 9–15 nm region, plus narrower windows around 19.2 nm, 25.5 nm, 28.4 nm, and 33.5 nm, fully covering the temperature range of ~2–25 MK.

Because full-Sun irradiances of lines formed at lower temperatures (~\(<2\) MK, e.g., Fe ix 17.1) are typically only modestly increased from the flare and, often, are decreased from eruption-associated coronal dimming (e.g., Mason et al. 2014), we do not currently analyze these cooler lines. Spatially resolved studies (e.g., Inglis & Christe 2014) could help to better address this in the future.

4 http://www.lmsal.com/solarsoft/
3. RESULTS AND DISCUSSION

For a first look at real data, we chose two GOES X-class flares—which have been shown likely to achieve the hottest temperatures (Caspì et al. 2014)—jointly observed by EVE and RHESSI: 2011 February 15 (X2.2) and 2011 March 9 (X1.5). The impulsive and early decay phase of each flare was treated as described above: subdivided into ∼60 s intervals, preprocessed, then fit with a DEM (and the additional model components for RHESSI).

Figure 3 shows the detailed results of one run of the model fit for two time periods, during the rise and HXR peak, of the February 15 event. Overall, the model spectra fit the observations well; there are systematics in the normalized residuals for both instruments, but the results appear reasonable. Although the $\chi^2$ is high, we feel it is acceptable given that we are fitting large regions of the spectra subject to unknown systematic uncertainties. For EVE, the residuals are largest around the brightest lines, and may be due to assuming a constant 0.1 nm FWHM resolution—the actual resolution varies slightly with wavelength—with no line broadening. Additionally, chromospheric and transition region lines (∼0.01–1 MK) are not modeled but may be blended with coronal lines; the most notable example is He II 25.6 nm, clearly not well-fit in the model and partly blended with Fe xxiv 25.5 nm. Higher-order contributions from the MEGS-A grating are also possibly present in the observed spectra, but are not modeled. For RHESSI, the residuals appear largest around the Fe/Fe–Ni line complexes, possibly due to insensitivity of the fit to these model parameters as they affect only a few out of many hundreds of data points. Despite the adequate intercalibration (Figure 1), as we are combining data from two unique instruments, inaccuracies in the respective instrument calibrations would contribute to the combined residuals, although this is difficult to quantify. Ionization nonequilibrium is not likely to be a significant contributor to the residuals in either instrument, as the high densities (∼10^{11–12} cm$^{-3}$) typically observed in X-class flares (e.g., Caspi & Lin 2010; Caspi et al. 2014) suggest ionization timescales for Fe of ∼10 s (∼1 s for the hotter ions, xvi–xxv; cf. Jordan 1970), much shorter than our ∼60 s integration times.

Figure 4 shows the time-series average DEM fits for both events. There is a clear progression from broader, hotter DEMs earlier (rise phase and peak) to narrower, cooler DEMs later (decay phase). Conspicuously absent, however, is any significant, distinct super-hot ($T_e \gtrsim 30$ MK; log $T_e \gtrsim 7.5$—cf. Lin et al. 1981) emission at any time during either flare. (The possible super-hot component during the earliest intervals of the March 9 event is weakly constrained due to poor statistics, and the RHESSI HXR spectra during these intervals are equally well-fit by non-thermal power-laws.) While a bimodal DEM, as has been inferred by previous studies (e.g., Jakimiec et al. 1988; Longcope et al. 2010; Caspi & Lin 2010; Warren et al. 2013; J. M. McTiernan & A. Caspi 2014, in preparation), is present during many intervals, the higher-temperature peak is typically at only ∼12–20 MK, well below the super-hot temperatures that would be expected from X-class events (Caspì et al. 2014), although we note that these two flares are fairly low X-class.
The peaks or plateaus in the joint-instrument DEM are typically at lower temperatures than the dual-isothermal approximations obtained from the RHESSI-only pre-fits, which is not unexpected since the continuous DEM predicts additional HXR flux from temperatures between and below the isothermally approximated values, requiring less high-temperature emission in the DEM to explain the remaining flux observed by RHESSI.

Because our model includes a non-thermal component, we are able to obtain limits on the non-thermal low-energy cutoff (Holman et al. 2011). Importantly, the EVE data provide a lower bound for the cutoff energy, unobtainable from RHESSI alone, as pushing the cutoff too low reduces the thermal model below what is supported by the EVE spectra; this is currently under study.

4. SUMMARY AND CONCLUSIONS

SDO/EVE’s unprecedented solar EUV observations, coupled with RHESSI’s groundbreaking X-ray measurements, provide an ideal toolkit to study the evolution of both thermal plasma and accelerated, non-thermal electrons in solar flares.

We have developed a technique to simultaneously analyze the EVE and X-ray spectra from EVE and RHESSI, respectively, enabling derivation of the most accurate, self-consistent flare DEMs to date covering the entire coronal temperature range observed in flares, ~2 to ~50 MK. We have validated this algorithm on artificial test data, and applied it to two X-class flares from the current solar cycle. The derived time-series DEMs progress clearly from broad, hot distributions to narrower, cooler distributions, and a bimodal signature is often observed. The DEMs suggest that fitting RHESSI data in isolation, particularly with isothermals, may bias the results towards higher temperatures (cf. Ryan et al. 2014), highlighting the need for the additional constraints provided by EVE. Indeed, these DEMs show little evidence of a significant super-hot component, as would be naively expected from X-class flares based on the results of Caspi et al. (2014). However, we note that most of the RHESSI dual-isothermal pre-fits also yielded no super-hot component, the only exceptions being the earliest one or two intervals of each event, which also had the broadest and highest-temperature DEMs; these two flares may simply be cooler events. Further, despite their X-class intensity, both flares were highly impulsive, short-duration events with little or no HXR emission ~100 keV; such flares may experience different heating than the eruptive, long-duration events more typically seen with these intensities (such as 10 of the 12 X-class flares in Caspi et al. 2014). The 2011 August 9 X6.9 flare, a longer-duration eruptive event, is currently under study and will help address this question.

One essential application of this DEM analysis will be to constrain the non-thermal electron low-energy cutoff. Comparing the thermal and non-thermal electron populations, and their evolution, would help to test different theories of plasma heating or to evaluate the partition between different heating mechanisms (e.g., Liu et al. 2013; Oreshina & Oreshina 2013).

Our technique also has important geospace applications, in helping determine the solar SXR loading on Earth’s upper atmosphere. The wavelength range of ~0.4–5 nm is thought to contain the bulk of flare-radiated energy (Rodgers et al. 2006) and is the primary driver of ionospheric dynamics. Very few spectrally resolved observations exist in this range, and the detailed dynamics depend critically on the altitude of energy deposition, strongly determined by the (unknown) spectral energy distribution (e.g., Sojka et al. 2013). Our physically modeled DEMs allow generation of synthetic X-ray spectra in this unobserved passband, filling a crucial gap in the needed input for modeling of upper atmospheric behavior.

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CASI, McTIERNAN, & WARREN

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