Front-end electronics and data acquisition system for a multi-wire 3D gas tracker

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Abstract

This paper presents the design and implementation of the front-end electronics and the data acquisition (DAQ) system for readout of multi-wire drift chambers (MWDC). Apart from the conventional drift time measurement the system delivers the hit position along the wire utilizing the charge division technique. The system consists of preamplifiers, and analog and digital boards sending data to a back-end computer via an Ethernet interface. The data logging software formats the received data and enables an easy access to the data analysis software. The use of specially designed preamplifiers and peak detectors allows the charge-division readout of the low resistance signal wire. The implication of the charge-division circuitry onto the drift time measurement was studied and the overall performance of the electronic system was evaluated in dedicated off-line tests.

Keywords: data acquisition system, hardware, data logging software
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1. Introduction

Precision measurements in neutron and nuclear decay offer a sensitive window to search for new physics beyond the standard electroweak model and allow also the determination of the fundamental weak vector coupling. Recent analyses based on the effective field theory performed in e.g. [1, 2] show that in processes involving the lightest quarks the neutron and nuclear decay will compete with experiments at highest energy accelerators. For instance, data taken at the LHC is currently probing these interactions at the $10^{-2}$ level (relative to the standard weak interactions), with the potential to reach the $10^{-3}$ level. In some of the $\beta$ decay correlation measurements there are prospects to reach experimental sensitivities between $10^{-3}$ and $10^{-4}$ making these observables interesting probes for searches of new physics originating at TeV scale. The most direct access to the exotic tensor interaction in $\beta$ decay is to measure the Fierz term (coefficient $b$) or the beta-neutrino correlation coefficient $a$ in a pure Gamow-Teller transition [3]. The $b$ coefficient shows up as a tiny energy dependent ($1/E$) departure of the $\beta$ spectrum from its V-A (standard model) shape. The smallness of the potential $b$ contribution requires that other corrections to the spectrum shape of the same order are included in the analysis. Indeed, according to [4, 5] the recoil terms also affect the spectrum shape with their main contribution being proportional to $E$. In order to disentangle these effects the detector efficiency for $\beta$ particle as a function of energy must be known with the precision better than $10^{-3}$ [6]. The dominating contribution in the systematic uncertainty comes from back-scattering and out-scattering of electrons from the detector. Monte Carlo simulation of this effect is helpful, however, it introduces its own uncertainty as the input parameters are known with limited accuracy. Monte Carlo simulation would reflect the real situation better after it is adjusted to real experimental data of a particular measurement setup.

The described in this article electronics was designed for a spectrometer capable of direct registration of the back-scattering events, thus providing reference data for the Monte Carlo calculation of the detector efficiency. The spectrometer itself is still in an R&D phase undergoing detailed tests and tuning. It will be a subject of a separate paper together with the performance benchmark [7]. In this paper, its concept will be described only in a minimum extent at the beginning of Section 2 to explain the requirements imposed onto the front-end
electronics and DAQ. The rest of Section 2 is devoted to the electronic system architecture. The test results were obtained with a help of a signal generator and are presented in Section 3. Therein the resulting time spectrum and the charge asymmetry distribution representing the typical performance are shown. The asymmetry spectra were obtained using a dedicated tester to simulate the hit position on the wire with adjustable resistance division (potentiometer). Conducting the electronic benchmark tests without a detector was chosen by purpose. The detector itself is still not fully understood. Therefore it was important to assess the purely electronic contribution to the performance parameters of the spectrometer. The paper ends with a short summary and outlook for the future experiment.

2. System architecture

In order to facilitate the identification of the electrons impinging onto and scattered from the energy detector (e.g. Si detector, scintillator) a low-Z and low-mass tracker must be applied. One of the attractive options is a low pressure multi-wire drift chamber with minimum number of necessary wires in order to maximize the detector transparency. This condition can be fulfilled by a hexagonal wire geometry and the charge division technique allowing for a 3D track reconstruction without major distortion of the electron energy measurement. The hexagonal wire geometry is not the only one considered in the project. The rectangular (planar) wire configuration is the next suitable alternative. The multi-wire drift chamber is based on the small prototype described in Ref. [8] and will be operated with He/isobutan gas mixture (ranging from 70%/30% to 90%/10%) at lowered pressure (down to 300 mbar). It consists of 10 sense wire planes (8 wires for each plane) separated with 24 field wire planes forming the double F-F-S-F-F-S-F-F-S-F-F-S-F-F-S-F-F structure (F- denotes a field wire plane, S- denotes a signal wire plane). The distance between neighboring signal planes is 15 mm, with wires within a plane being separated by 17.32 mm. This wire plane structure leads to the hexagonal cell geometry as shown in Fig. [1]. Each cell consists of a very thin anode wire (NiCr alloy, 25 µm diameter) with resistance of about 20 ohms/cm surrounded by 6 cathode field wires forming a hexagonal cuboid. All wires are soldered to pads of the printed circuit board (PCB) frames. The chamber is equipped with a two-dimensional positioning system for a beta source installed in the central region of the detector, between both parts of the MWDC structure. In the initial configuration, the electrons detected in two plastic scintillators installed at both sides of the MWDC provide the time reference signal for the drift time measurement. The PMT signals are also used as a trigger for the MWDC and electron energy detector readout. Acquiring the drift time and the pulse height asymmetry at both ends of the responding signal wire one can establish the electron path across the chamber cell. The system provides charge-division position sensing in the direction parallel to the wires as well as precise drift time measurement. The expected position resolution across wires is limited by angular straggling of primary electrons travelling in the gas and accounts to about 450 µm as shown in Ref. [8]. In this situation, the precision of the drift time measurement of about 200 ps is more than needed as it corresponds to about 50 µm for the operating conditions in view. The MWDC will be operated in homogeneous magnetic field oriented parallel to wires providing a rough electron transverse momentum filter. The position information along the wires will be used for the identification of the sequence of the cells passed by the electrons and for distinguishing between the electrons impinging onto and scattered from the energy detector. This is why the modest position resolution of a few mm is sufficient for that direction.

Crucial in the design is the analog part of the system. In principle, commercial digitizers (TDC, ADC) could be applied at the end of the acquisition chain. However, it has been decided to equip the data channels with custom digital boards with adjusted specifications such that the whole system is handy, cost effective and scalable. The applied on board processing (FPGA) allows for ac-
quiring up to 15 000 triggered events per second which gives a comfortable factor of 10 reserve as compared to the application in view.

The described modular electronic system consists of three main parts: (i) preamplifiers, (ii) analog cards containing the peak detector and constant fraction discriminator (CFD), and (iii) the digital boards containing analog to digital converters, ADC and TDC. The data logging software on the PC constitutes the data receiver. The signal from each signal plane is processed by means of one electronics module, which consists of two analog cards and one digital board providing 16 ADC and 8 TDC channels. A corresponding block diagram is shown in Fig. 2. The signals from both ends of a signal wire are fed to inputs of preamplifiers located directly on the detector frames (see Fig. 3) in order to minimize the input noise and protect the signal from EM interferences. The signals from the preamplifiers are received by the analog cards which drive the ADC inputs and produce the TDC STOP signals. The digital data is transmitted via a LAN port to the back-end computer where the process of receiving, sorting and formatting to a complete physical event is accomplished by the data logging software. The card configuration settings and control is done via a RS485 port. Detailed description of each part of the system is presented in the following sections.

2.1. Preamplifiers

Signal readout from both ends of the relatively low resistance wire requires application of dedicated preamplifiers. Low resistance wires force the use of fast preamplifiers with low input impedance since the change of signal amplitudes and thus position resolution depends on the input impedance. The lower the input impedance, the higher the voltage difference that is registered at the wire ends. The preamplifiers work in current-mode with the input impedance below 5 ohms and the input stage bandwidth exceeding 300 MHz. The preamplifier circuit is shown in Fig. 4. The first stage of the preamplifier is a current-voltage converter with two fast bipolar transistors. The voltage pulse is fed to the high-pass filter which cuts off the DC bias and the slow varying components of the signal. The first filter is followed by a low-pass filter integrating the pulse with 100 ns time constant. The signal is then amplified and transmitted to the differential amplifier used to drive the transmission line. Differential signal transmission increases the external noise immunity when unshielded twisted pair ribbon cables are used to connect the preamplifier outputs to the inputs of the analog module. An example of input and output signals from a preamplifier is shown in Fig. 5. The input current signal was reproduced with the help of Qucs package [9]. Using bipolar transistors increases the preamplifier immunity to possible discharges in the gas chamber. Adding more robust input protection (e.g. a series resistor and protection diodes) was abandoned as it would increase the input noise and impedance thus
worsening the charge division resolution. However, accidental discharges cannot be avoided completely and the need for replacement of damaged preamplifiers was taken into account in the design. The individual preamplifiers are arranged as small piggyback PCB cards mounted directly on the wire frames (Fig. 5). They can be replaced easily without unplugging cables or any other major intervention in the setup. Additionally, such a solution minimizes the length of the unshielded connection between the wire end and the preamplifier input.

2.2. Analog cards

The block diagram of the analog circuit is presented in Fig. 6.

Differential signals from the preamplifiers are processed by the analog circuit. The signals from both wire ends are split into two branches. In the first branch, signals from both ends are added and fed to the CFD which produces a TTL pulse for the time-to-digital converter (TDC). A delay correction prior to the summing is not needed unless it is of the order of 1 ns (corresponding to about 20 cm cable length) significantly less than the rise time of real signals from the MWDC. Care is taken to assure the equal length of the connecting cables. If the sum of the analog signals is higher than the set value of the CFD threshold, the STOP signal for the TDC is generated. The schematic of the CFD is presented in Fig. 7.

Fig. 5: (color on-line) Upper panel: Preamplifier input current signal reproduced with the help of Qucs package. Lower panel: Measured output signal from the preamplifier. Blue and red lines denote positive and negative output signals.

Fig. 6: Block diagram of the analog board.

Fig. 7: Schematic of the Constant Fraction Discriminator.

In the second branch, both signals are transmitted to the fast peak-hold detectors, which are responsible for detection of the pulse amplitudes in a given gate time. Peak-hold detectors are used to stretch the pulse to the length acceptable for the ADC converters. The schematic of the peak-hold detector is shown in Fig. 8.

The gate and hold signals are produced by a programmable timing circuit located on the analog board which uses the TDC START signal (external trigger) as a time reference. The analog boards hosting the analog signal processing circuits are equipped with built-in controllers for setting the thresholds of the CFD as well.
as the gate and the hold timing. The set values are fed to the controller over a slow-control bus (RS485).

2.3. Digital cards

The block diagram of the digital circuit is presented in Fig. 9. The digital board consists of 16 channels of 14-bit ADCs (AD8099 chips made by Analog Devices) and a 8 channel TDC (TDC-GPX chip made by Acam Messelectronic Gmbh). Using 14-bit ADCs was dictated by the large dynamic range of the amplitudes of the signals from a gas detector operated in the proportional mode. The high bit ADC resolution increases the overall performance of the charge division method. The time resolution of the used TDC chip is significantly better than actually needed (83 ps binning resolution with the typical standard deviation of 77 ps, from the device datasheet). The board measures the delay between the common START (trigger) pulse and the individual STOP signals received from the constant fraction discriminators, as well as the amplitudes of the pulses delivered by the peak-hold detectors. The STOP signals are delayed by a 50 ns delay line in order to compensate for the delays introduced in the analog circuit generating the START signal. The ADCs sample the signals with 20 Msp sampling rate. The FPGA chip (Xilinx Spartan XC3S400) receives the data from the ADCs and the TDC, buffers them, provides zero-suppression and time-stamping, and formats the data frame as described in the next paragraph. The flow diagram of the acquisition algorithm and the data transmission algorithm are shown in Fig. 10. The data frame is transmitted to the integrated Ethernet/UDP stack (Wiznet W5100) which is responsible for communication with the back-end PC hosting the acquisition and control software. The digital board is also equipped with a built-in controller providing initialization and configuration of ADC, TDC and Ethernet chips. This controller is accessible via the slow-control bus.
2.4. Data formatter

The data are transferred to the computer in a form of UDP frames. The data formatter creates a frame for each recorded event. The frame length is a multiple of 96 bytes. The first part of the frame is a fixed header used for synchronization. Next follows the time tag (32 bit unsigned integer value) which allows the time synchronization between multiple modules. The third part of the frame consists of 16 ADC conversion values (16 bit unsigned integer values). The fourth part contains the TDC values - up to eight 24-bit unsigned integer values. Only the non-zero values of the TDC data are transmitted. The time tagging is used because the Ethernet protocol does not ensure the sequence of the data packet delivery. The time-tags are shared between modules and used by the data-logging software to restore the correct data order. Fig. 11 shows the structure of the data frame.

2.5. Slow-control bus and interface

The slow-control interface allows setting the CFD thresholds and timing configuration parameters of the module. Each board in the module is equipped with its own controller connected to the slow-control bus. Individual addresses of the controllers are set with dip-switches installed on the boards. The controllers use a simple ASCII protocol via RS485 interface. In order to connect the bus to the acquisition computer a commercial RS485-USB converter is used. An example of the initialization data record necessary for establishing the communication over the slow-control bus is presented on Fig. 12. The slow-control parameter set is read back every second and refreshed on a computer display (Fig. 13).

2.6. Data logging software

The data logging software provides the parallel reception of data frames sent by the modules and the event building i.e. collection of data generated by the system upon a single hardware trigger signal and uniquely identified by a time tag. Each module is identified by its IP address and uses a specific UDP port corresponding to this address. The architecture of the logging software is shown in Fig. 14. The software is multi-threaded and consists of the main thread (the parent) that is responsible for sorting and recording events, and the child
processes whose task is to receive data from the acquisition modules and to send them via data streams to the main thread. Upon initialization the program reads the configuration file with the addresses of the modules and opens sockets for all modules to communicate with them. The sockets listen to the different UDP ports correlated with the module address. The main loop creates a set of pipes (data streams) and receives the data structures from the child processes. The size of the frame is checked and followed by the pre-selection of cases. If the TDC value is greater than zero, the event is pushed to the FIFO queue of the child process. The event is then read from the FIFO by the parent process in which the data structures from all child processes are sorted (using the quick-sort algorithm) by the time tags and formatted as one complete event. This event will be identified with a particle crossing the chamber planes and directed for physics analysis. The received events are accessible for on-line histogramming analysis by means of a shared memory mechanism and simultaneously saved on a hard disk for later off-line analysis.

The cases with empty TDC conversion value (STOP signal below CFD threshold) are used to calculate the baseline offset for the ADC. In this method it is assumed that the lack of the TDC signal in a pair of channels (serving a particular sense wire) means that the cell did not fire so the corresponding ADCs deliver the offset values. This in-flight offset calculation method can be disabled and the fixed offsets can be introduced with the help of the slow-control interface. The baseline offsets subtracted (in flight) from the conversion value are appended to the data structures for each event so that the subtraction is reversible. Finally, the data structure consists of a time tag, module number, channel number, ADC1 value, ADC2 value, TDC value, ADC1 offset and ADC2 offset, respectively. It is worth mentioning that the amplitude thresholds which are needed for the histogram building were offset from the ADC baseline values such that only the electronic noise was eliminated. The ultimate adjustment of the thresholds was postponed to the detector efficiency tuning phase. It will depend on the gas mixture composition and pressure as well as on the operation voltage.

3. Test results

The described DAQ system was tested using a special tester simulating the wire chamber signals. The wire itself is replaced by a potentiometer with the range reflecting the real wire resistance (see Fig. 15). The tester utilizes two preamplifiers and signal cables connected to the selected module. A 400 mV high and 200 ns long simulator input is synchronized with a TTL signal triggering the system. The simulator input was fed by a triangle signal with 10 ns rise time and 200 ns fall time adjusted to the expected detector pulse shape obtained from the GARFIELD (Ref. [10]) simulation. The hit position corresponds to a unique potentiometer setting. The corresponding ADC asymmetry distributions are shown in Fig. 16. The zero point defines the middle of the wire. The difference between the individual asymmetry spectra reflects the varying relation between the signal and noise amplitudes. In Fig. 17 sample results from one channel are presented.

Fig. 18 shows a number of centroids the ADC pulse height asymmetry distributions

\[
A_{ADC} = V_A - V_B \\
A_{ADC} = V_A + V_B
\]

(1)

acquired at different potentiometer asymmetry settings

\[
A_{pot} = R_A - R_B \\
A_{pot} = R_A + R_B
\]

(2)
Fig. 16: (color on-line) ADC asymmetry spectra taken for 12 different potentiometer settings. The red hashed peak corresponds to $A_{pot} = 0$.

Mean 7049
RMS 204.3

Counts
0 5000 1000

Mean 7049
RMS 201.8

Fig. 17: Sample results for a single channel of one acquisition module. The two top plots present collected ADC values corresponding to both the upper part and the lower part of the wire. The two lower plots show the ADC asymmetry and converted TDC values in nanoseconds.

where $V_A$ and $V_B$ are the ADC1, ADC2 amplitudes and $R_A, R_B$ are the resistances for selected potentiometer settings. The results were obtained for varying resistance division corresponding to the charge collected at both wire ends. The obtained ADC asymmetry is a monotonic function of the resistance asymmetry. The centroids of these distributions are drawn with $1\sigma$ error bars in Fig. 18. The polynomial function fitted to the centroids represents the position calibration. Drawing the error band allows the extraction of the position resolution as shown in Fig. 19. The uncertainty of the relative position resolution is connected with the interpolation procedure. It shows that the charge division method determines the position with a resolution between 1.2 and 3.6 mm for a 240 mm long wire. When feeding a given preamplifier pair with the simulator signals the noise distributions on the neighboring channels were acquired as well. Neither a baseline change nor an increase of noise were observed, meaning that the electronic crosstalk is negligible. This does not assure there

Fig. 18: ADC pulse height asymmetry (centroids of the peaks from Fig. 16) as a function of the resistance asymmetry $A_{pot}$ (Eqn. 2). The insert is a zoomed part of the graph showing the error bars equal to $\pm 1\sigma$ of the peak distributions plotted in Fig. 16. Dotted lines interpolate the error bar ends.

Fig. 19: Position resolution $\Delta L / L_0$ expressed in units of length of the wire deduced from Fig. 18. A 500 ohm wire resistance was assumed which corresponds to a 25 $\mu$m NiCr wire of $L_0 = 240$ mm length. The uncertainty is connected with the interpolation procedure (see text).
is no crosstalk when the electronic system is attached to the gas detector. Inductive crosstalk between the wires depends on the configuration and on the operating conditions and is beyond the scope of this paper.

In the second test, the response of the TDC measurement was investigated as a function of the input signal asymmetry. Fig. 20 shows a typical example. The impact of the charge division on the TDC measurement is found to be less than 1 ns in the entire charge asymmetry range. For the application in view, a 1 ns drift time corresponds to about 250 µm distance (from GARFIELD simulations).

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Fig. 20: TDC values as a function of the ADC asymmetry. The error bars correspond to the standard deviation of the TDC histogram.

4. Conclusions and outlook

The front-end electronics and DAQ system described here was designed to be used with a special multi-wire drift chamber for tracking low energy electrons from nuclear β decays. It incorporates both the drift time and charge division measurements allowing for efficient 3D determination of the electron tracks with very few and only parallel sense wires. The performed tests show that the electronic contribution to the drift time measurement uncertainty is less than 1 ns corresponding to about 250 µm position uncertainty at the expected electron drift velocities. Such result is satisfactory since in the planned experiment the track position resolution will be dominated by the electron angular straggling effects in the gas as shown in Monte Carlo simulations and confirmed in the small prototype test described in Ref. [1]. The necessary spatial resolution of the electron track position determined from the drift time need not to be better than 500 µm. The uncertainty of the track position obtained from the charge division measurement varies from 0.5% at wire ends to 1.5% in the center corresponding to 1.2 and 3.6 mm, respectively, for 24 cm long wires (25 µm NiCr). This result is sufficient for the application in view: identification of the cell sequence passed by an electron spiraling in an axial magnetic field and resolving possible double track ambiguities appearing in the 2D projection.

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