GEODESICS IN THE HEISENBERG GROUP $\mathbb{H}^n$ BY WAY OF FOURIER SERIES

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Abstract. We provide a new and elementary proof for the structure of geodesics in the Heisenberg group $\mathbb{H}^n$. The proof is based on an adaptation of Hurwitz's proof of the isoperimetric inequality in $\mathbb{R}^2$ using Fourier series. We also prove that the Carnot-Carathéodory metric is real analytic away from the center of the group.

1. Introduction

The Heisenberg group $\mathbb{H}^n$ is $\mathbb{C}^n \times \mathbb{R} = \mathbb{R}^{2n+1}$ given the structure of a Lie group with multiplication

$$(z, t) \ast (z', t') = \left( z + z', t + t' + 2\text{Im} \sum_{j=1}^{n} z_j z'_j \right)$$

$$= \left( x_1 + x'_1, \ldots, x_n + x'_n, y_1 + y'_1, \ldots, y_n + y'_n, t + t' + 2 \sum_{j=1}^{n} (x'_j y_j - x_j y'_j) \right)$$

with Lie algebra $\mathfrak{g}$ whose basis of left invariant vector fields at any point $p = (x_1, \ldots, x_n, y_1, \ldots, y_n, t) \in \mathbb{H}^n$ is

$$X_j(p) = \frac{\partial}{\partial x_j} + 2y_j \frac{\partial}{\partial t}, \quad Y_j(p) = \frac{\partial}{\partial y_j} - 2x_j \frac{\partial}{\partial t}, \quad T = \frac{\partial}{\partial t}, \quad j = 1, 2, \ldots, n.$$ We call $H\mathbb{H}^n = \text{span}\{X_1, \ldots, X_n, Y_1, \ldots, Y_n\}$ the horizontal distribution on $\mathbb{H}^n$, and denote by $H_p\mathbb{H}^n$ the horizontal space at $p$. An absolutely continuous curve $\Gamma : [0, S] \to \mathbb{R}^{2n+1}$ is said to be horizontal if $\dot{\Gamma}(s) \in H_{\Gamma(s)}\mathbb{H}^n$ for almost every $s \in [0, S]$. It is easy to see that the horizontal distribution is the kernel of the standard contact form

$$\alpha = dt + 2 \sum_{j=1}^{n} (x_j dy_j - y_j dx_j).$$ That is, $H_p\mathbb{H}^n = \ker \alpha(p)$. Hence it follows that an absolutely continuous curve $\Gamma = (x_1, \ldots, x_n, y_1, \ldots, y_n, t)$ is horizontal if and only if

$$t(s) = 2 \sum_{j=1}^{n} (\dot{x}_j(s)y_j(s) - x_j(s)\dot{y}_j(s))$$

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for almost every $s \in [0, S]$. This means that

$$t(s) - t(0) = 2 \sum_{j=1}^{n} \int_{0}^{s} (\dot{x}_j(\tau)y_j(\tau) - x_j(\tau)\dot{y}_j(\tau)) \, d\tau$$

for every $s \in [0, S]$. Suppose $\gamma = (x_1, \ldots, x_n, y_1, \ldots, y_n) : [0, S] \to \mathbb{R}^{2n}$ is absolutely continuous. If a value for $t(0)$ is fixed, then (1.1) gives a unique horizontal curve $\Gamma = (\gamma(s), t(s))$ whose projection onto the first $2n$ coordinates equals $\gamma$. We call this curve $\Gamma$ a horizontal lift of $\gamma$.

If the curve $\gamma : [0, S] \to \mathbb{R}^{2n}$ is closed and $\Gamma = (\gamma, t)$ is a horizontal lift of $\gamma$, then it follows from Green’s Theorem that the total change in height $t(S) - t(0)$ equals $-4$ times the sum of the signed areas enclosed by the projections $\gamma_j$ of the curve to each $x_jy_j$-plane. Since the curves in the $x_jy_j$-planes may have self-intersections, the signed area has to take multiplicity of the components of the complement of $\gamma_j$ into account. This multiplicity can be defined in terms of the winding number. We will not provide more details here as this interpretation of the change of the height will not play any role in our argument and it is mentioned here only to give an additional vantage point to the geometry of the problem.

We equip the horizontal distribution $H\mathbb{H}^n$ with a left invariant Riemannian metric so that the vector fields $X_j$ and $Y_j$ are orthonormal at every point $p \in \mathbb{H}^n$. Note that this metric is only defined on $H\mathbb{H}^n$ and not on $T\mathbb{H}^n$. If

$$v = \sum_{j=1}^{n} \left( a_j \frac{\partial}{\partial x_j} \bigg|_p + b_j \frac{\partial}{\partial y_j} \bigg|_p \right) + c \frac{\partial}{\partial t} \bigg|_p \in H_p\mathbb{H}^n,$$

then it is easy to see that

$$v = \sum_{j=1}^{n} a_j X_j(p) + b_j Y_j(p)$$

and hence $c = 2 \sum_{j=1}^{n} (a_j y_j(p) - b_j x_j(p))$. Clearly $|v|_H \leq |v|_E$, where

$$|v|_H = \sqrt{\sum_{j=1}^{n} a_j^2 + b_j^2} \quad \text{and} \quad |v|_E = \sqrt{\left( \sum_{j=1}^{n} a_j^2 + b_j^2 \right) + c^2}$$

are the lengths of $v$ with respect to the metric in $H_p\mathbb{H}^n$ and the Euclidean metric in $\mathbb{R}^{2n+1}$ respectively. On compact sets in $\mathbb{R}^{2n+1}$ the coefficient $c$ is bounded by $C|v|_H$, so for any compact set $K \subset \mathbb{R}^{2n+1}$ there is a constant $C(K) \geq 1$ such that

$$|v|_H \leq |v|_E \leq C(K)|v|_H \quad \text{for all} \; p \in K \text{ and } v \in H_p\mathbb{H}^n.$$

A horizontal curve $\Gamma : [a, b] \to \mathbb{H}^n$ satisfies

$$\dot{\Gamma}(s) = \sum_{j=1}^{n} \dot{x}_j(s)X_j(\Gamma(s)) + \dot{y}_j(s)Y_j(\Gamma(s)).$$

Its length with respect to the metric in $H\mathbb{H}^n$ equals

$$\ell_H(\Gamma) := \int_{a}^{b} |\dot{\Gamma}(s)|_H \, ds = \int_{a}^{b} \sqrt{\sum_{j=1}^{n} (\dot{x}_j(s))^2 + (\dot{y}_j(s))^2} \, ds.$$
The Carnot-Carathéodory metric $d_{cc}$ in $\mathbb{H}^n$ is defined as the infimum of lengths $\ell_H(\Gamma)$ of horizontal curves connecting given two points. It is well known that any two points can be connected by a horizontal curve (we will actually prove it) and hence $d_{cc}$ is a true metric. It follows from (1.2) that for any compact set $K \subset \mathbb{R}^{2n+1}$ there is a constant $C = C(K)$ such that

\begin{equation}
|p - q| \leq C d_{cc}(p, q) \quad \text{for all } p, q \in K.
\end{equation}

If $\Gamma : [a, b] \to X$ is a (continuous) curve in a metric space $(X, d)$, then the length of $\Gamma$ is defined by

$$
\ell_d(\Gamma) = \sup \sum_{i=0}^{n-1} d(\Gamma(s_i), \Gamma(s_{i+1})),
$$

where the supremum is taken over all $n \in \mathbb{N}$ and all partitions $a = s_0 \leq s_1 \leq \ldots \leq s_n = b$. In particular if $\Gamma : [a, b] \to \mathbb{H}^n$ is a curve in the Heisenberg group, then its length with respect to the Carnot-Carathéodory metric is

\begin{equation}
\ell_{cc}(\Gamma) = \sup \sum_{i=0}^{n-1} d_{cc}(\Gamma(s_i), \Gamma(s_{i+1})).
\end{equation}

It follows immediately from the definition that if $\Gamma$ is a horizontal curve, then $\ell_{cc}(\Gamma) \leq \ell_H(\Gamma)$; hence every horizontal curve is rectifiable (i.e. of finite length), but it is not obvious, whether $\ell_{cc}(\Gamma) = \ell_H(\Gamma)$. It is also not clear whether every rectifiable curve can be reparametrized as a horizontal curve. Actually all this is true and we have

**Proposition 1.1.** In $\mathbb{H}^n$ we have

1. Any horizontal curve $\Gamma$ is rectifiable and $\ell_{cc}(\Gamma) = \ell_H(\Gamma)$.
2. Lipschitz curves are horizontal.
3. Every rectifiable curve admits a 1-Lipschitz parametrization (and hence it is horizontal with respect to this parametrization).

Thus up to a reparametrization, the class of rectifiable curves coincides with the class of horizontal ones. Proposition 1.1 is well known, but it is difficult to find a good reference that would provide a straightforward proof. For the sake of completeness, we decided to provide a proof of this result in the Appendix. For a proof of (1) in more generality, see Theorem 1.3.5 in [15].

A geodesic from $p$ to $q$ is a curve of shortest length between the two points. It is already well known that any two points in the Heisenberg group can be connected by a geodesic. Also, the structure of every geodesic in a form of an explicit parameterization is known. The proofs in the case of $\mathbb{H}^1$ can be found in [2, 5, 8, 14] and the general case of $\mathbb{H}^n$ is treated in [1, 3, 16].

Notice that the length $\ell_H(\Gamma)$ of a horizontal curve $\Gamma = (\gamma, t)$ given by (1.3) equals the usual Euclidean length $\ell_E(\gamma)$ in $\mathbb{R}^{2n}$ of the projection $\gamma$.

If $n = 1$, the structure of geodesics can be obtained via the two dimensional isoperimetric inequality (see [2, 5, 14]). Consider a horizontal curve $\Gamma = (\gamma, t)$ in $\mathbb{H}^1$ connecting the origin to some point $q = (0, 0, T)$ with $T \neq 0$. The length of $\Gamma$ equals the length of its projection $\gamma$ to $\mathbb{R}^2$ (which is a closed curve). Also, by (1.1), the change $T$ in the height of $\Gamma$ must
equal $-4$ times the signed area enclosed by $\gamma$. Thus the projection of any horizontal curve connecting 0 to $q$ must enclose the same area $|T|/4$, and finding a geodesic which connects 0 to $q$ reduces to a problem of finding a shortest closed curve $\gamma$ enclosing a fixed area. Thus the isoperimetric inequality implies that $\Gamma$ will have smallest length when $\gamma$ is a circle. Then the $t$ component of $\Gamma$ is determined by (1.1) and one obtains an explicit parametrization of the geodesics in $\mathbb{H}^1$ connecting the origin to a point on the $t$-axis. Such geodesics pass through all points $(x_0, y_0, t_0)$, $t_0 \neq 0$ in $\mathbb{H}^1$. If $q = (x_0, y_0, 0)$, then it is easy to see that the segment $0q$ connecting the origin to $q$ is a geodesic. This describes all geodesics connecting the origin to any other point in $\mathbb{H}^1$. Due to the left-invariance of the vector fields $X$ and $Y$, parameterizations for geodesics between arbitrary points in $\mathbb{H}^1$ may be found by left multiplication of the geodesics discussed above.

This elegant argument however, does not apply to $\mathbb{H}^n$ when $n > 1$ and known proofs of the structure of geodesics in $\mathbb{H}^n$ are based on the Pontryagin maximum principle, [1,3,16]. The purpose of this note is to provide a straightforward and elementary argument leading to an explicit parameterization of geodesics in $\mathbb{H}^n$ (Theorem 2.1). Our argument is based on Hurwitz’s proof [11], of the isoperimetric inequality in $\mathbb{R}^2$ involving Fourier series. See [17] for a similar argument in a different but related setting. As an application of our method we also prove that the Carnot-Carathéodory metric is real analytic away from the center of the group, see Theorem 3.1. This improves a result of Monti [15,16] who proved that this distance is $C^\infty$ smooth away from the center. We also find a formula for the Carnot-Carathéodory distance (Corollary 3.2) that can be used to compute the terms of the power series expansion explicitly.

The paper is organized as follows. In Section 2 we will state and prove the result about the structure of the geodesics in $\mathbb{H}^n$, and in Section 3 we use this structure to show that the distance function in $\mathbb{H}^n$ is analytic away from the center. In Section 4 we address the problem of comparing different geodesics (there are infinitely many of them) connecting the origin to a point $(0,0,T) \in \mathbb{H}^n$ on the $t$-axis. Finally, in the Appendix we prove Proposition 1.1.

We hope that the paper will be of interest to those who are new to the subject of the Heisenberg group and for that reason we decided to make the paper self-contained.

2. The structure of geodesics

Any horizontal curve $\Gamma$ is rectifiable and using arc-length parametrization we can assume that the speed $|\dot{\Gamma}|_H$ of $\Gamma : [0, \ell_H(\Gamma)] \to \mathbb{H}^n$ equals 1. Then we can reparametrize it as a curve of constant speed defined on $[0,1]$, so we can assume that $\Gamma : [0,1] \to \mathbb{H}^n$ satisfies

\begin{equation}
\sum_{j=1}^n \dot{x}_j(s)^2 + \dot{y}_j(s)^2 = \ell_H(\Gamma)^2 = \ell_{cc}(\Gamma)^2 \quad \text{for almost all } s \in [0,1].
\end{equation}

On the other hand any rectifiable curve in $\mathbb{H}^n$ can be reparametrized as a horizontal curve via the arc length parameterization (Proposition 1.1) and hence when looking for length minimizing curves (geodesics) it suffices to restrict to horizontal curves $\Gamma : [0,1] \to \mathbb{H}^n$ satisfying (2.1).
Since the left translation in $\mathbb{H}^n$ is an isometry, it suffices to investigate geodesics connecting the origin $0 \in \mathbb{H}^n$ to another point in $\mathbb{H}^n$. Indeed, if $\Gamma$ is a geodesic connecting $0$ to $p^{-1} \ast q$, then $p \ast \Gamma$ is a geodesic connecting $p$ to $q$.

If $q$ belongs to the subspace $\mathbb{R}^{2n} \times \{0\} \subset \mathbb{R}^{2n+1} = \mathbb{H}^n$, then it is easy to check that the straight line $\Gamma(s) = sq$, $s \in [0, 1]$ is a unique geodesic (up to a reparametrization) connecting $0$ to $q$. Indeed, it is easy to check that $\Gamma$ is horizontal and its length $\ell_{cc}(\Gamma) = \ell_H(\Gamma)$ equals the Euclidean length $|0q|$ of the segment $0q$, because $\Gamma$ is equal to its projection $\gamma$. For any other horizontal curve $\hat{\Gamma} = (\hat{\gamma}, \hat{t})$ connecting $0$ to $q$ the projection $\hat{\gamma}$ on $\mathbb{R}^{2n}$ would not be a segment (since horizontal lifts of curves are unique up to vertical shifts) and hence we would have $\ell_{cc}(\hat{\Gamma}) = \ell_H(\hat{\Gamma}) = \ell_E(\hat{\gamma}) > |0q| = \ell_{cc}(\Gamma)$ which proves that $\hat{\Gamma}$ cannot be a geodesic.

In Theorem 2.1 we will describe the structure of geodesics in $\mathbb{H}^n$ connecting the origin to a point $(0, 0, T) \in \mathbb{R}^{2n} \times \mathbb{R} = \mathbb{H}^n$, $T \neq 0$, lying on the $t$-axis. Later we will see (Corollary 2.3) that these curves describe all geodesics in $\mathbb{H}^n$ connecting $0$ to $q \notin \mathbb{R}^{2n} \times \{0\}$. The geodesics connecting $0$ to $q \in \mathbb{R}^{2n} \times \{0\}$ have been described above.

**Theorem 2.1.** A horizontal curve

$$\Gamma(s) = (x(s), y(s), t(s)) = (x_1(s), \ldots, x_n(s), y_1(s), \ldots, y_n(s), t(s)) : [0, 1] \to \mathbb{H}^n$$

of constant speed, connecting the origin $\Gamma(0) = (0, 0, 0) \in \mathbb{R}^{2n} \times \mathbb{R} = \mathbb{H}^n$ to a point $\Gamma(1) = (0, 0, \pm T)$, $T > 0$, on the $t$-axis is a geodesic if and only if

$$x_j(s) = A_j(1 - \cos(2\pi s)) \mp B_j \sin(2\pi s)$$

$$y_j(s) = B_j(1 - \cos(2\pi s)) \pm A_j \sin(2\pi s)$$

for $j = 1, 2, \ldots, n$ and

$$t(s) = \pm T \left( s - \frac{\sin(2\pi s)}{2\pi} \right)$$

where $A_1, \ldots, A_n, B_1, \ldots, B_n$ are any real numbers such that $4\pi \sum_{j=1}^n (A_j^2 + B_j^2) = T$.

Observe that if $\Gamma(1) = (0, 0, +T)$ the equations (2.2) give a constant-speed parametrization of negatively oriented circles in each of the $x_jy_j$-planes, centered at $(A_j, B_j)$, and of radius $\sqrt{A_j^2 + B_j^2}$. Each circle passes through the origin at $s = 0$. The signed area of such a circle equals $-\pi(A_j^2 + B_j^2)$. Thus the change in height $t(1) - t(0)$ which is $-4$ times the sum of the signed areas of the projections of the curve on the $x_jy_j$-planes equals

$$(-4) \sum_{j=1}^n (-\pi(A_j^2 + B_j^2)) = T.$$ 

Clearly this must be the case, because $\Gamma$ connects the origin to $(0, 0, T)$. Any collection of circles in the $x_jy_j$-planes passing through the origin and having radii $r_j \geq 0$ are projections of a geodesic connecting the origin to the point $(0, 0, T)$ where $T = 4\pi \sum_{j=1}^n r_j^2$. In particular we can find a geodesic for which only one projection is a nontrivial circle (all other radii are zero) and a geodesic for which all projections are non-trivial circles. That suggests that the geodesics connecting $(0, 0, 0)$ to $(0, 0, T)$ may have many different shapes. This is, however, an incorrect intuition. As we will see in Section 4 all such geodesics are obtained from one through a rotation of $\mathbb{R}^{2n+1}$ about the $t$ axis. This rotation is also an isometric
mapping of \( \mathbb{H}^n \). The above reasoning applies also to the case when \( \Gamma(1) = (0, 0, -T) \) with the only difference that now the circles are positively oriented.

**Proof of Theorem 2.1.** Suppose first that \( \Gamma : [0, 1] \to \mathbb{H}^n \) is any horizontal curve of constant speed connecting the origin to the point \((0, 0, +T), T > 0 \). Recall from (2.1) that

\[
\sum_{j=1}^{n} \dot{x}_j(s)^2 + \dot{y}_j(s)^2 = \ell_{cc}(\Gamma)^2 =: L^2.
\]

In particular the functions \( x_j \) and \( y_j \) are \( L \)-Lipschitz continuous and \( x_j(0) = y_j(0) = x_j(1) = y_j(1) = 0 \). Hence the functions \( x_j, y_j \) extend to \( 1 \)-periodic Lipschitz functions on \( \mathbb{R} \) and so we can use Fourier series to investigate them. We will follow notation used in [6]. By Parseval’s identity,

\[
L^2 = \sum_{j=1}^{n} \int_{0}^{1} |\dot{x}_j(s)|^2 + |\dot{y}_j(s)|^2 \, ds = \sum_{j=1}^{n} \sum_{k \in \mathbb{Z}} |\dot{x}_j(k)|^2 + |\dot{y}_j(k)|^2
\]

(2.3)

\[
= \sum_{j=1}^{n} \sum_{k \in \mathbb{Z}} 4\pi^2 k^2 (|\dot{x}_j(k)|^2 + |\dot{y}_j(k)|^2).
\]

Also, we have by (1.1)

\[
T = 2 \sum_{j=1}^{n} \int_{0}^{1} (\dot{x}_j(s)y_j(s) - \dot{y}_j(s)x_j(s)) \, ds.
\]

Thus by noting that \( \dot{x}_j(s) = \overline{\dot{x}_j(s)} \) and \( \dot{y}_j(s) = \overline{\dot{y}_j(s)} \), we may apply Parseval’s theorem to this pair of inner products and find

\[
T = 2 \sum_{j=1}^{n} \left( \sum_{k \in \mathbb{Z}} \overline{\dot{x}_j(k)}\dot{y}_j(k) - \sum_{k \in \mathbb{Z}} \overline{\dot{y}_j(k)}\dot{x}_j(k) \right)
\]

\[
= 2 \sum_{j=1}^{n} \sum_{k \in \mathbb{Z}} 2\pi ki \left( \overline{\dot{y}_j(k)}\dot{x}_j(k) - \overline{\dot{x}_j(k)}\dot{y}_j(k) \right)
\]

(2.4)

\[
= 4\pi \sum_{j=1}^{n} \sum_{k \in \mathbb{Z}} k \cdot 2 \text{Im}(\overline{\dot{x}_j(k)}\dot{y}_j(k)),
\]

because \( i(\bar{a}b - \bar{a}b) = 2 \text{Im}(\bar{a}b) \). Subtracting (2.4) from (2.3) gives

\[
\frac{L^2}{4\pi^2} - \frac{T}{4\pi} = \sum_{j=1}^{n} \left[ \sum_{k \in \mathbb{Z}} k^2 \left( |\dot{x}_j(k)|^2 + |\dot{y}_j(k)|^2 \right) - k \cdot 2 \text{Im}(\overline{\dot{x}_j(k)}\dot{y}_j(k)) \right]
\]

\[
= \sum_{j=1}^{n} \left[ \sum_{k \in \mathbb{Z}} (k^2 - |k|) \left( |\dot{x}_j(k)|^2 + |\dot{y}_j(k)|^2 \right) \right.
\]

\[
+ |k| \left( |\dot{x}_j(k)|^2 - 2 \text{sgn}(k) \text{Im}(\overline{\dot{x}_j(k)}\dot{y}_j(k)) + |\dot{y}_j(k)|^2 \right) \left. \right]
\]

(2.5)

\[
= \sum_{j=1}^{n} \left[ \sum_{k \in \mathbb{Z}} (k^2 - |k|) \left( |\dot{x}_j(k)|^2 + |\dot{y}_j(k)|^2 \right) + \sum_{k \in \mathbb{Z}} |k||\dot{x}_j(k) + i \text{sgn}(k)\dot{y}_j(k)|^2 \right].
\]
The last equality follows from the identity $|a + bi|^2 = |a|^2 - 2 \text{Im}(\bar{a}b) + |b|^2$ which holds for all $a, b \in \mathbb{C}$. Since every term in this last sum is non-negative, it follows that $\frac{L^2}{4\pi^2} - \frac{T}{4\pi} \geq 0$. Thus, we have $L^2 \geq \pi T$. Now $\Gamma$ is a geodesic if and only if $L^2 = \pi T$ which happens if and only if each of the two sums in (\ref{eq:geodesic}) equals zero. Since $k^2 - |k| > 0$ for $|k| \geq 2$, the first of the two sums vanishes if and only if $\hat{x}_j(k) = \hat{y}_j(k) = 0$ for every $|k| \geq 2$ and $j = 1, 2, \ldots, n$. Hence nontrivial terms in the last sum correspond to $k = \pm 1$ and the last sum vanishes if and only if $\hat{x}_j(\pm 1) = -i \text{sgn}(\pm 1) \hat{y}_j(\pm 1)$. That is, for every $j = 1, \ldots, n$,

\begin{equation}
\hat{x}_j(\pm 1) = -i \hat{y}_j(1) \quad \text{and} \quad \hat{x}_j(-1) = i \hat{y}_j(-1).
\end{equation}

Now since each $x_j$ and $y_j$ is Lipschitz, their Fourier series converge uniformly on $[0, 1]$. Note that the only non-zero terms in the Fourier series appear when $|k| \leq 1$. Thus $\Gamma$ is a geodesic if and only if (\ref{eq:geodesic}) is satisfied and for every $s \in [0, 1]$ and $j = 1, \ldots, n$

\begin{align}
x_j(s) &= \hat{x}_j(-1)e^{-2\pi is} + \hat{x}_j(0) + \hat{x}_j(1)e^{2\pi is} \\
y_j(s) &= \hat{y}_j(-1)e^{-2\pi is} + \hat{y}_j(0) + \hat{y}_j(1)e^{2\pi is}.
\end{align}

In particular, $0 = x_j(0) = \hat{x}_j(-1) + \hat{x}_j(0) + \hat{x}_j(1)$ and hence $\hat{x}_j(0) = -\hat{x}_j(-1) - \hat{x}_j(1)$ for each $j = 1, \ldots, n$. This together with Euler’s formula gives

\begin{align*}
x_j(s) &= \hat{x}_j(-1)(e^{-2\pi is} - 1) + \hat{x}_j(1)(e^{2\pi is} - 1) \\
&= -(\hat{x}_j(-1) + \hat{x}_j(1))(1 - \cos(2\pi s)) + (-i\hat{x}_j(-1) + i\hat{x}_j(1))\sin(2\pi s) \\
&= -(\hat{x}_j(-1) + \hat{x}_j(1))(1 - \cos(2\pi s)) + (\hat{y}_j(-1) + \hat{y}_j(1))\sin(2\pi s).
\end{align*}

The last equality follows from (\ref{eq:geodesic}). Similarly, we have

\begin{align*}
y_j(s) &= -(\hat{y}_j(-1) + \hat{y}_j(1))(1 - \cos(2\pi s)) - (\hat{x}_j(-1) + \hat{x}_j(1))\sin(2\pi s).
\end{align*}

If we write $A_j = -(\hat{x}_j(-1) + \hat{x}_j(1))$ and $B_j = -(\hat{y}_j(-1) + \hat{y}_j(1))$, then we have the desired parameterization

\begin{align}
x_j(s) &= A_j(1 - \cos(2\pi s)) - B_j\sin(2\pi s) \\
y_j(s) &= B_j(1 - \cos(2\pi s)) + A_j\sin(2\pi s)
\end{align}

which is the $(0, 0, +T)$ case of (\ref{eq:geodesic}). Note that it follows directly from the definition of Fourier coefficients that the numbers $A_j$, $B_j$ are real. Now the formula for the $t$ component of $\Gamma$ follows from (\ref{eq:geodesic}); the integral is easy to compute due to numerous cancellations.

Suppose now that $\Gamma: [0, 1] \to \mathbb{H}^n$ is any horizontal curve of constant speed connecting the origin to the point $(0, 0, -T)$, $T > 0$. Then $\Gamma = (x(s), y(s), t(s))$ is a geodesic if and only if

\[ \Gamma(s) = (\tilde{x}(s), \tilde{y}(s), \tilde{t}(s)) = (x(1 - s), y(1 - s), t(1 - s) + T) \]

is a geodesic connecting $\tilde{\Gamma}(0) = (0, 0, 0)$ and $\tilde{\Gamma}(1) = (0, 0, T)$ because reversing parametrization does not change the length of a curve and the mapping $(x, y, t) \mapsto (x, y, t + T)$ (vertical lift by $T$) is an isometry on $\mathbb{H}^n$. Therefore $\tilde{\Gamma} = (\tilde{x}, \tilde{y}, \tilde{t})$ must have the form (\ref{eq:geodesic}). Hence the $(0, 0, -T)$ case of (\ref{eq:geodesic}) follows from (\ref{eq:geodesic}) by replacing $s$ with $1 - s$.

\begin{remark}
Notice that the inequality $L^2 \geq \pi T$ acts as a sort of isoperimetric inequality for curves in $\mathbb{R}^{2n}$. Since the change in height $T$ of the geodesic $\Gamma$ is related to the sum $A$ of the signed areas of the projected curves via the relationship $T = -4A$ as seen in (\ref{eq:geodesic}), our inequality becomes $L^2 \geq 4\pi |A|$. In the case $n = 1$, this is the classical isoperimetric inequality known since antiquity.
\end{remark}
To have a better understanding of the structure of geodesics it is convenient to use complex notation following the identification of \( \mathbb{R}^{2n} \) with \( \mathbb{C}^n \) given by
\[
\mathbb{R}^{2n} \ni (x, y) = (x_1, \ldots, x_n, y_1, \ldots, y_n) \leftrightarrow (x_1 + iy_1, \ldots, x_n + iy_n) = x + iy \in \mathbb{C}^n.
\]
With this notation the geodesics from Theorem 2.1 connecting the origin to \((0, 0, \pm T)\), \(T > 0\) can be represented as
\[
(2.9) \quad \Gamma(s) = \left((1 - e^{\mp 2\pi is})(A + iB), t(s)\right)
\]
where \(A = (A_1, \ldots, A_n)\), \(B = (B_1, \ldots, B_n)\) are such that \(4\pi |A + iB|^2 = T\) and
\[
t(s) = \pm T \left(s - \frac{\sin(2\pi s)}{2\pi}\right).
\]

Theorem 2.1 and a discussion preceding it describes geodesics connecting the origin to points either on the \(t\)-axis \((0, 0, \pm T)\), \(T > 0\) or in \(\mathbb{R}^{2n} \times \{0\}\). The question now is how to describe geodesics connecting the origin to a point \(q\) which is neither on the \(t\)-axis nor in \(\mathbb{R}^{2n} \times \{0\}\). It turns out that geodesics described in Theorem 2.1 cover the entire space \(\mathbb{H}^n \setminus (\mathbb{R}^{2n} \times \{0\})\) and we have

**Corollary 2.3.** For any \(q \in \mathbb{H}^n\) which is neither in the \(t\)-axis nor in the subspace \(\mathbb{R}^{2n} \times \{0\}\) there is a unique geodesic connecting the origin to \(q\). This geodesic is a part of a geodesic connecting the origin to a point on the \(t\)-axis.

**Proof.** Let \(q = (c_1, \ldots, c_n, d_1, \ldots, d_n, h)\) be such that \(h \neq 0\) and \(c_j, d_j\) are not all zero. We can write \(q = (c + id, h) \in \mathbb{C}^n \times \mathbb{R}\). First we will construct a geodesic \(\Gamma\) given by (2.9) so that \(\Gamma(s_0) = q\) for some \(s_0 \in (0, 1)\). Clearly such a geodesic is a part of a geodesic connecting the origin to a point on the \(t\)-axis. Then we will prove that this is a unique geodesic (up to a reparametrization) connecting the origin to \(q\). Assume that \(h > 0\) (the case \(h < 0\) is similar). We will find a geodesic passing through \(q\) that connects \((0, 0, 0)\) to \((0, 0, T)\), for some \(T > 0\). (If \(h < 0\) we find \(\Gamma\) that connects \((0, 0, 0)\) to \((0, 0, -T)\).) It suffices to show that there is a point \(A + iB \in \mathbb{C}^n\) such that the system of equations
\[
(2.10) \quad (1 - e^{-2\pi is})(A + iB) = c + id, \quad 4\pi |A + iB|^2 \left(s - \frac{\sin(2\pi s)}{2\pi}\right) = h
\]
has a solution \(s_0 \in (0, 1)\). We have \(A + iB = (c + id)/(1 - e^{-2\pi is})\) and hence
\[
(2.11) \quad 2\pi \frac{|c + id|^2}{1 - \cos(2\pi s)} \left(s - \frac{\sin(2\pi s)}{2\pi}\right) = h.
\]
This equation has a unique solution \(s_0 \in (0, 1)\) because the function on the left hand side is an increasing diffeomorphism of \((0, 1)\) onto \((0, \infty)\). We proved that, among geodesics connecting \((0, 0, 0)\) to points \((0, 0, T)\), there is a unique geodesic \(\Gamma_q\) passing through \(q\). Suppose now that \(\tilde{\Gamma}\) is any geodesic connecting \((0, 0, 0)\) to \(q\). Gluing \(\tilde{\Gamma}\) with \(\Gamma_q|_{[s_0, 1]}\) we obtain a geodesic connecting \((0, 0, 0)\) to \((0, 0, T)\) and hence (perhaps after a reparametrization) it must coincide with \(\Gamma_q\). This proves uniqueness. \(\square\)

We will now use the proof of Corollary 2.3 to find a formula for the Carnot-Carathéodory distance between \((0, 0)\) and \(q = (z, h), z \neq 0, h > 0\). We will need this formula in the next
section. Let
\begin{equation}
H(s) = \frac{2\pi s - \sin(2\pi s)}{1 - \cos(2\pi s)} : (0, 1) \to (0, \infty)
\end{equation}
be the diffeomorphism of $(0, 1)$ onto $(0, \infty)$ described in (2.11). Let
\begin{equation}
\Gamma(s) = ((1 - e^{-2\pi is})(A + iB), t(s))
\end{equation}
be the geodesic from the proof of Corollary 2.3 that passes through $q$ at $s_0 \in (0, 1)$. We proved that $s_0$ is a solution to (2.11) and hence $s_0$ is a function of $q$ given by
\begin{equation}
\label{eqn:2.12}
s_0(q) = H^{-1}(h|z|^{-2}).
\end{equation}
Note that $A + iB = z/(1 - e^{-2\pi is_0})$, so
\begin{equation}
|A + iB| = \frac{|z|}{\sqrt{2(1 - \cos(2\pi s_0))}}.
\end{equation}
Hence
\begin{equation}
\sqrt{\sum_{j=1}^{n} \dot{x}_j^2(s) + \dot{y}_j^2(s)} = L = \sqrt{\pi T} = 2\pi |A + iB| = \frac{2\pi |z|}{\sqrt{2(1 - \cos(2\pi s_0))}}.
\end{equation}
Where $L$ is the length of $\Gamma$, and $T$ satisfies $\Gamma(1) = (0, 0, T)$. Thus
\begin{equation}
d_{cc}(0, q) = \int_{0}^{s_0} \sqrt{\sum_{j=1}^{n} \dot{x}_j^2(s) + \dot{y}_j^2(s)} = \frac{2\pi s_0 |z|}{\sqrt{2(1 - \cos(2\pi s_0))}}.
\end{equation}

3. Analyticity of the Carnot-Carathéodory metric

The center of the Heisenberg group $\mathbb{H}^n$ is $Z = \{(z, h) \in \mathbb{H}^n \mid z = 0\}$. It is well known that the distance function in $\mathbb{H}^n$ is $C^\infty$ smooth away from the center \cite{15, 16}, but, via the above arguments, we will now see that this distance function is in fact real analytic.

**Theorem 3.1.** The Carnot-Carathéodory distance $d_{cc} : \mathbb{R}^{2n+1} \times \mathbb{R}^{2n+1} \to \mathbb{R}$ is real analytic on the set
\begin{equation}
\{(p, q) \in \mathbb{H}^n \times \mathbb{H}^n = \mathbb{R}^{2n+1} \times \mathbb{R}^{2n+1} : q^{-1} \ast p \notin Z\}.
\end{equation}

**Proof.** In the proof we will make a frequent use of a well known fact that a composition of real analytic functions is real analytic, \cite[Proposition 2.2.8]{13}. It suffices to prove that the function $d_0(p) = d_{cc}(0, p)$ is real analytic on $\mathbb{H}^n \setminus Z$. Indeed, $w(p, q) = q^{-1} \ast p$ is real analytic as it is a polynomial. Also, $d_{cc}(p, q) = (d_0 \circ w)(p, q)$, so real analyticity of $d_0$ on $\mathbb{H}^n \setminus Z$ will imply that $d_{cc}$ is real analytic on $w^{-1}(\mathbb{H}^n \setminus Z) = \{(p, q) \in \mathbb{H}^n \times \mathbb{H}^n : q^{-1} \ast p \notin Z\}$.

Define $H : (-1, 1) \to \mathbb{R}$ as
\begin{equation}
H(s) = \frac{2\pi s - \sin(2\pi s)}{1 - \cos(2\pi s)} = \frac{2\pi s}{3!} - \frac{(2\pi s)^3}{5!} + \frac{(2\pi s)^5}{7!} - \ldots.
\end{equation}
Here, we divided by a common factor of $(2\pi s)^2$ in the two power series on the right hand side. That is, the denominator equals $(1 + \cos(2\pi s))(2\pi s)^{-2}$ which does not vanish on $(-1, 1)$. This implies that $H$ is real analytic on $(-1, 1)$. Indeed, considering $s$ as a complex
variable, we see that \( H(s) \) is holomorphic (and hence analytic) in an open set containing \((-1, 1)\) as a ratio of two holomorphic functions with non-vanishing denominator.

As we pointed out in (2.12), the function \( H \) is an increasing diffeomorphism of \((0, 1)\) onto \((0, \infty)\). Since it is odd and \( H'(0) = 2\pi/3 \neq 0 \), \( H \) is a real analytic diffeomorphism of \((-1, 1)\) onto \(\mathbb{R} \). Again, using a holomorphic function argument we see that \( H^{-1} : \mathbb{R} \to (-1, 1) \) is real analytic.

The function \( z \mapsto |z|^{-2} \) is analytic on \(\mathbb{R}^{2n} \setminus \{0\} \) (as a composition of a polynomial \( z \mapsto |z|^2 \) and an analytic function \(1/x\)), so the function \((z, h) \mapsto h|z|^{-2}\) is analytic in \(\mathbb{H}^n \setminus Z \). Hence also \( s_0(q) = H^{-1}(h|z|^{-2}) \) is analytic on \(\mathbb{H}^n \setminus Z \).

Fix \( q = (z, h) \in \mathbb{H}^n \setminus Z \) with \( h > 0 \). Then by (2.13)

\[
(3.2) \quad d_0(q) = \frac{2\pi s_0|z|}{\sqrt{2(1 - \cos(2\pi s_0))}}.
\]

Since \( H(s_0) = h|z|^{-2} \), formula (3.1) yields

\[
2\pi s_0 = (1 - \cos(2\pi s_0))h|z|^{-2} + \sin(2\pi s_0).
\]

Substituting \(2\pi s_0 \) in the numerator of the right hand side of (3.2) gives

\[
d_0(q) = \frac{h\sqrt{1 - \cos(2\pi s_0)}}{\sqrt{2}|z|} + \frac{|z|\sin(2\pi s_0)}{\sqrt{2}\sqrt{1 - \cos(2\pi s_0)}} = h\sin(\pi s_0)|z|^{-1} + |z|\cos(\pi s_0)
\]

where we used the trigonometric identities

\[
\sqrt{1 - \cos(2\pi s_0)} = \sqrt{2|\sin(\pi s_0)|} = \sqrt{2}\sin(\pi s_0) \quad \text{and} \quad \frac{\sin(2\pi s_0)}{\sin(\pi s_0)} = 2\cos(\pi s_0).
\]

To treat the case \( h \leq 0 \) we will define \( s_0(q) = H^{-1}(h|z|^{-2}) \) for any \( q = (z, h), z \neq 0 \). Previously we defined \( s_0(q) \) only when \( z \neq 0 \) and \( h > 0 \). It is easy to check that the mapping

\[
q = (x, y, t) \mapsto (z, t) \mapsto \bar{q} = (\bar{z}, -t) = (x, -y, -t)
\]

is an isometry of the Heisenberg group, so \( d_0(q) = d_0(\bar{q}) \).

If \( h < 0 \) and \( \bar{q} = (\bar{z}, -h) \), then

\[
s_0(q) = H^{-1}(h|z|^{-2}) = -H^{-1}(-h|\bar{z}|^{-2}) = -s_0(q)
\]

and hence

\[
d_0(q) = d_0(\bar{q}) = -h\sin(\pi s_0(\bar{q}))|\bar{z}|^{-1} + |\bar{z}|\cos(\pi s_0(\bar{q})) = h\sin(\pi s_0(q))|z|^{-1} + |z|\cos(\pi s_0(q)).
\]

In the case \( h = 0 \), \( \Gamma \) is a straight line in \(\mathbb{R}^{2n} \) from the origin to \( q \), and so \( d_0(q) = |z| \).

Therefore

\[
d_0(q) = h\sin(\pi s_0(q))|z|^{-1} + |z|\cos(\pi s_0(q))
\]

for every \( q = (z, h) \in \mathbb{H}^n \setminus Z \), and so \( d_0 \) is analytic on \(\mathbb{H}^n \setminus Z \). \(\square\)

We also proved
Corollary 3.2. The Carnot-Carathéodory distance between the origin \((0,0)\) and \((z,h)\), \(z \neq 0\) equals

\[d_{cc}((0,0),(z,h)) = h \sin(\pi H^{-1}(h|z|^2))|z|^{-1} + |z| \cos(\pi H^{-1}(h|z|^2)).\]

One can easily use this result to compute the terms of the power series expansion of \((z,h) \mapsto d_{cc}((0,0),(z,h))\), \(z \neq 0\).

4. Classification of non-unique geodesics

Any point \((0,0,\pm T)\), \(T > 0\) on the \(t\) axis can be connected to the origin by infinitely many geodesics. The purpose of this section is to show that actually all such geodesics are obtained from one geodesic by a linear mapping which fixes the \(t\)-axis. This map is an isometry of \(\mathbb{H}^n\) and also an isometry of \(\mathbb{R}^{2n+1}\).

Proposition 4.1. If \(\Gamma_1 : [0,1] \to \mathbb{H}^n\) and \(\Gamma_2 : [0,1] \to \mathbb{H}^n\) are constant-speed geodesics with \(\Gamma_1(0) = \Gamma_2(0) = (0,0,0)\) and \(\Gamma_1(1) = \Gamma_2(1) = (0,0,\pm T)\) with \(T > 0\), then we can write \(\Gamma_2 = V \circ \Gamma_1\) where \(V\) is a linear isometry in \(\mathbb{H}^n\) which fixes the \(t\)-coordinate. The map \(V\) is also an isometry of \(\mathbb{R}^{2n+1}\), a rotation about the \(t\)-axis.

Proof. Consider geodesics \(\Gamma_1 = (\gamma_1, t)\) and \(\Gamma_2 = (\gamma_2, t)\) defined in the statement of the proposition. As in the discussion before Corollary 2.3 we consider \(\gamma_1\) and \(\gamma_2\) as functions into \(\mathbb{C}^n\) rather than into \(\mathbb{R}^{2n}\) and write

\[\gamma_1(s) = (1 - e^{\mp 2\pi is})(A + iB), \quad \gamma_2(s) = (1 - e^{\mp 2\pi is})(C + iD),\]

where \(4\pi|A+iB|^2 = 4\pi|C+iD|^2 = T\). We claim that there is a unitary matrix \(U \in U(n, \mathbb{C})\) such that \(U(A+iB) = C+iD\). Indeed, for any \(0 \neq z \in \mathbb{C}^n\), use the Gram-Schmidt process to extend \(\{z/|z|\}\) to an orthonormal basis of \(\mathbb{C}^n\) and define \(W_z\) to be the matrix whose columns are these basis vectors. Here, we consider orthogonality with respect to the standard Hermitian inner product \(\langle u, v \rangle = \sum_{j=1}^n u_j \overline{v_j}\). Then \(W_z \in U(n, \mathbb{C})\) and \(W_z e_1 = z/|z|\) where \(\{e_1, \ldots, e_n\}\) is the standard basis of \(\mathbb{C}^n\). Thus the desired operator is \(U = W_{C+iD} \circ W_{A+iB}^{-1}\).

Define the linear map \(V : \mathbb{C}^n \times \mathbb{R} \to \mathbb{C}^n \times \mathbb{R}\) by \(V(z,t) = (Uz,t)\). Since

\[U \left((1 - e^{\mp 2\pi is})(A + iB)\right) = (1 - e^{\mp 2\pi is})(C + iD),\]

for every \(s \in [0,1]\) and since \(V\) fixes the \(t\)-component of \(\mathbb{C}^n \times \mathbb{R}\), we have \(V \circ \Gamma_1 = \Gamma_2\).

We now prove that \(V\) is an isometry on \(\mathbb{H}^n\). Indeed, suppose \(p,q \in \mathbb{H}^n\) and \(\Gamma = (\gamma, t) : [0,1] \to \mathbb{H}^n\) is a geodesic connecting them. Then \(V \circ \Gamma = (U \circ \gamma, t)\). Since \(\Gamma\) is horizontal, it is easy to check that \(i(s) = 2 \text{Im} \langle \gamma(s), \dot{\gamma}(s) \rangle_{\mathbb{C}}\) for almost every \(s \in [0,1]\). Unitary operators preserve the standard inner product on \(\mathbb{C}^n\), and so

\[i(s) = 2 \text{Im} \langle \gamma(s), \dot{\gamma}(s) \rangle_{\mathbb{C}} = 2 \text{Im} \langle (U \circ \gamma)(s), (U \circ \dot{\gamma})(s) \rangle_{\mathbb{C}} = 2 \text{Im} \left(\frac{d}{ds}((U \circ \gamma)(s), \frac{d}{ds}(U \circ \gamma)(s))_{\mathbb{C}}\right),\]
for almost every $s \in [0,1]$. That is, $V \circ \Gamma$ is horizontal. Also,

$$\ell_H(\Gamma) = \int_0^1 \sqrt{\langle \dot{\gamma}(s), \dot{\gamma}(s) \rangle_C} \, ds = \int_0^1 \sqrt{\langle (U \circ \dot{\gamma})(s), (U \circ \dot{\gamma})(s) \rangle_C} \, ds = \ell_H(V \circ \Gamma).$$

Thus $d_{cc}(Vp, Vq) \leq \ell_H(\Gamma) = d_{cc}(p, q)$. Since $U$ is invertible and $U^{-1} \in U(n, \mathbb{C})$, we may argue similarly to show that $d_{cc}(p, q) = d_{cc}(V^{-1}Vp, V^{-1}Vq) \leq d_{cc}(Vp, Vq)$, and so $V$ is an isometry on $\mathbb{H}^n$. Clearly unitary transformations of $\mathbb{C}^n$ are also orientation preserving isometries of $\mathbb{R}^{2n}$ and hence $V$ is a rotation of $\mathbb{R}^{2n+1}$ about the $t$-axis. 

5. Appendix

Proof of Proposition 5.1 (1) The components of any horizontal curve $\Gamma$ are absolutely continuous, and so their derivatives are integrable. Thus the inequality $\ell_{cc}(\Gamma) \leq \ell_H(\Gamma) < \infty$ yields rectifiability of horizontal curves. It remains to prove that $\ell_{cc}(\Gamma) \geq \ell_H(\Gamma)$. Extend the Riemannian tensor defined on the horizontal distribution $H\mathbb{H}^n$ to a Riemannian tensor $g$ in $\mathbb{R}^{2n+1}$. For example we may require that the vector fields $X_j, Y_j, T$ are orthonormal at every point of $\mathbb{R}^{2n+1}$. The Riemannian tensor $g$ defines a metric $d_g$ in $\mathbb{R}^{2n+1}$ in a standard way as the infimum of lengths of curves connecting two given points, where the length of an absolutely continuous curve $\alpha : [a, b] \to \mathbb{R}^{2n+1}$ is defined as the integral

$$\ell_g(\alpha) = \int_a^b \sqrt{g(\dot{\alpha}(s), \dot{\alpha}(s))} \, ds.$$

This is the same approach that was used to define the Carnot-Carathéodory metric. For horizontal curves $\Gamma$, we have $\ell_g(\Gamma) = \ell_H(\Gamma)$, and so it is obvious that $d_g(p, q) \leq d_{cc}(p, q)$ since we now take an infimum over a larger class of curves. It is a well known fact in Riemannian geometry that for an absolutely continuous curve $\alpha : [a, b] \to \mathbb{R}^{2n+1}$

$$\ell_g(\alpha) = \sum_{i=0}^{n-1} d_g(\alpha(s_i), \alpha(s_{i+1})),\]$$

where the supremum is taken over all $n \in \mathbb{N}$ and all partitions $a = s_0 \leq s_1 \leq \ldots \leq s_n = b$ as before. Hence if $\Gamma : [a, b] \to \mathbb{R}^{2n+1}$ is horizontal we have

$$\ell_H(\Gamma) = \ell_g(\Gamma) = \sup_{i=0}^{n-1} d_g(\Gamma(s_i), \Gamma(s_{i+1})) \leq \sup_{i=0}^{n-1} d_{cc}(\Gamma(s_i), \Gamma(s_{i+1})) = \ell_{cc}(\Gamma).$$

(2) If $\Gamma : [a, b] \to \mathbb{H}^n$ is Lipschitz, then by (1) it is also Lipschitz with respect to the Euclidean metric in $\mathbb{R}^{2n+1}$. Thus it is absolutely continuous and differentiable a.e. It remains to show that $\Gamma(s) \in H_{\Gamma(s)}\mathbb{H}^n$ a.e. We will actually show that this is true whenever $\Gamma$ is differentiable at $s$.

Suppose $\Gamma$ is differentiable at a point $s \in (a, b)$ and let $i_0 \in \mathbb{N}$ be such that $s + 2^{-i_0} \leq b$. Then $d_{cc}(\Gamma(s + 2^{-i}), \Gamma(s)) \leq L2^{-i}$ for some $L > 0$ and all $i \geq i_0$. By the definition of the Carnot-Carathéodory metric on $\mathbb{H}^n$, there is some horizontal curve $\eta^i : [0, 2^{-i}] \to \mathbb{H}^n$ which connects $\Gamma(s)$ to $\Gamma(s + 2^{-i})$ whose length approximates the distance between them. That is, we can choose a horizontal curve $\eta^i = (x^i, y^i, t^i)$ so that $\eta^i(0) = \Gamma(s)$, $\eta^i(2^{-i}) = \Gamma(s + 2^{-i})$.
and \( \ell_H(\eta^i) < 2L2^{-i} \). After a reparameterization, we may assume that \( \eta^i \) has constant speed \( |\dot{\eta}^i|_H < 2L \) on \([0, 2^{-i}]\). Since \( \eta^i \) is horizontal, we can write
\[
\dot{\eta}^i(\tau) = \sum_{j=1}^{n} \dot{x}^i_j(\tau)X_j(\eta^i(\tau)) + \dot{y}^i_j(\tau)Y_j(\eta^i(\tau))
\]
for almost every \( \tau \in [0, 2^{-i}] \). Now \( \Gamma(s+2^{-i}) - \Gamma(s) = \eta^i(2^{-i}) - \eta^i(0) = \int_0^{2^{-i}} \dot{\eta}^i(\tau) d\tau \), and so
\[
\frac{\Gamma(s+2^{-i}) - \Gamma(s)}{2^{-i}} = \frac{1}{2^{-i}} \int_0^{2^{-i}} \sum_{j=1}^{n} (\dot{x}^i_j(\tau)X_j(\eta^i(\tau)) + \dot{y}^i_j(\tau)Y_j(\eta^i(\tau))) d\tau
\]
(5.1)
\[
= \frac{1}{2^{-i}} \int_0^{2^{-i}} \sum_{j=1}^{n} \dot{x}^i_j(\tau)(X_j(\eta^i(\tau)) - X_j(\Gamma(s))) d\tau + \frac{1}{2^{-i}} \int_0^{2^{-i}} \sum_{j=1}^{n} \dot{y}^i_j(\tau)(Y_j(\eta^i(\tau)) - Y_j(\Gamma(s))) d\tau
\]
(5.2)
\[
+ \frac{1}{2^{-i}} \int_0^{2^{-i}} \sum_{j=1}^{n} (\dot{x}^i_j(\tau)X_j(\Gamma(s)) + \dot{y}^i_j(\tau)Y_j(\Gamma(s))) d\tau
\]

The images of all the curves \( \eta^i \) are contained in a compact subset of \( \mathbb{R}^{2n+1} \). Hence (1.2) yields
\[
\sup_{\tau \in [0, 2^{-i}]} |\eta^i(\tau) - \Gamma(s)| \leq \ell_E(\eta^i) \leq C(K)\ell_H(\eta^i) \leq C(K)L2^{-i+1} \to 0 \quad \text{as} \quad i \to \infty.
\]
Here \( \ell_E(\eta^i) \) stands for the Euclidean length of \( \eta^i \). Observe also that the functions \( \dot{x}^i_j \) and \( \dot{y}^i_j \) are bounded almost everywhere by the speed \( |\dot{\eta}^i|_H \) which is less than 2L. This and the continuity of the vector fields \( X_j, Y_j \) immediately imply that the sums (5.1) and (5.2) converge to zero as \( i \to \infty \). Also, by again using the uniform boundedness of the functions \( \dot{x}^i_j \) and \( \dot{y}^i_j \), we conclude that for some subsequence and all \( j = 1, 2, \ldots, n \) the averages \( 2^i \int_0^{2^{-i}} \dot{x}^i_j \) and \( 2^i \int_0^{2^{-i}} \dot{y}^i_j \) converge to some constants \( a_j \) and \( b_j \) respectively. Denote such a subsequence by \( i_k \). Therefore, since \( \hat{\Gamma}(s) \) exists we have
\[
\hat{\Gamma}(s) = \lim_{k \to \infty} \frac{\Gamma(s+2^{-i_k}) - \Gamma(s)}{2^{-i_k}} = \sum_{j=1}^{n} (a_jX_j(\Gamma(s)) + b_jY_j(\Gamma(s))) \in H_{\Gamma(s)}\mathbb{H}^n.
\]
This proves (2). Finally (3) follows from a general fact that a rectifiable curve in any metric space admits an arc-length parametrization with respect to which the curve is 1-Lipschitz, see e.g. [4] Proposition 2.5.9, [10] Theorem 3.2. □

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