Mid-rapidity charged hadron transverse spherocity in pp collisions simulated with Pythia

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Abstract

The pp collisions have been studied for a long time, however, there are still some effects which are not completely understood, such as the long range angular correlations and the flow patterns in high multiplicity events, which were recently discovered at the LHC. In a recent work it was demonstrated that in Pythia 8, multi-parton interactions and color reconnection can give some of the observed effects similar to the collective flow well known from heavy-ion collisions. Now using the same model, a study based on mid-rapidity charged hadron transverse spherocity is presented. The main purpose of this work is to show that a differential study combining multiplicity and event shapes opens the possibility to understand better the features of data, specially at high multiplicity.

Keywords: Multi-parton interactions, flow-like behavior, event shapes, pp collisions.

1. Introduction

Recent results at the LHC have uncovered hitherto unknown features in high multiplicity events. Perhaps the most interesting result is the discovery of sQGP-like (ridge-like and flow-like behavior) features in small systems like those created in pp and p-Pb collisions [1–6]. The origin of such effects is still far from being understood, for instance, it has been shown that hydro calculations, where the formation of a hot and dense QCD medium is implicitly assumed, can describe qualitatively many features of data [7–9]. But, other approaches suggest that the phenomenon can be produced by initial state effects [10–12]. For instance, Pythia 8.180 [13] gives an effect reminiscent of the collective flow well known from heavy-ion collisions [14]. This flow-like behavior is attributed to multi-parton interactions (MPI) and color reconnection (CR) [15].

The impression that the understanding of the initial stages of the heavy ion collisions is reasonably known contrasts with the opinion of a representative number of physicists who emphasize the fact that the understanding of processes transforming the initial stage into a hydrodynamical final state is still incomplete and requires further theoretical and experimental knowledge [16].

In this work the results of an event shape analysis, specifically, a selection on transverse spherocity of mid-rapidity charged hadrons, applied to events generated with Pythia 8.180 are presented. The goal of the study is to show that the use of event shapes allows to extract more information from data. Namely, this technique opens the possibility to isolate jetty-like (high $p_T$ jets) and isotropic (large number of low $Q^2$ partonic scatterings) events [17]. It is worth to notice that the procedure is inverse to the usual approach. Namely, in the majority of cases one tries to apply models, which successfully describe heavy ion collision data, to smaller system, sometimes forgetting that the premises valid in large systems are not at all satisfied in smaller ones. For example, the requirement that the Knudsen number
Figure 1. (Color online). Mid-rapidity ($|\eta| < 1$) charged hadron multiplicity as a function of the number of multi-parton interactions for pp collisions at $\sqrt{s} = 7$ TeV simulated with Pythia 8.180 tune 4C. Each multiplicity interval is normalized to one. The solid line illustrates the average number of MPI in a given multiplicity interval.

$K = \lambda/R$, where $\lambda$ is the mean free path and $R$ is the dimension of the system, must be small to allow for a rapid equilibration; is not really warranted in small systems or similarly, for particles that exhibit small interaction cross section with the medium.

With the exception of the anisotropic flow measurements [18], where the so-called non-flow contributions are well under control, most of the observables used to study the properties of the hot and dense QCD matter contain a mixture of contributions from the different components present in the collisions i.e., soft and hard QCD processes. For example, as illustrated in this work [19], the identified hadron production measured in central Pb-Pb collisions is completely different in the jet region and outside the jet peak (bulk). However, the low $p_T$ (< 3 GeV/c) part of the inclusive spectra is used to extract expansion velocity and the temperature at the kinetic freeze-out of the system [21]. In this context, it is argued that the implementation of an event shape analysis would allow for a better understanding of non-radial flow effects.

2. Transverse spherocity

In the present work, the mid-rapidity charged hadron transverse spherocity, $S_0$, is used to characterize the events through the geometrical distribution of the $p_T$'s of the charged hadrons, which is by definition collinear and infrared safe. The restriction to the transverse plane avoids the bias from the boost along the beam axis [22]. It is defined for a unit transverse vector $\hat{n}$ which minimizes the ratio below:

$$S_0 = \frac{\pi^2}{4} \left( \frac{\sum_i |p_T^i \times \hat{n}|}{\sum_i p_T^i} \right)^2. \quad (1)$$

By construction, the limits of the variable are related to specific configurations in the transverse plane

$$S_0 = \begin{cases} 
0 & \text{"pencil-like" limit (hard events)} \\
1 & \text{"isotropic" limit (soft events)} 
\end{cases}.$$ 

In this study, inelastic pp collisions were generated with Pythia 8.180 tune 4C, this tune describes qualitatively many features of the LHC pp data [23]. The event shape was computed considering only primary charged particles at mid-rapidity ($|\eta| < 1$) and in the transverse momentum interval $0.15 < p_T < 10$ GeV/c. Transverse spherocity was

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1. A similar result has been obtained from the analysis of p-Pb data using jet reconstruction and strange hadrons [20].
only defined for events with more than two hadrons. Different observables like jet production and identified hadron production were studied at mid-rapidity for different $S_0$ and multiplicity ($N_{ch}$) intervals.

Event shapes studied by experiments at LHC, e.g. ATLAS [24] and ALICE [25], have shown an interesting result: a good agreement between data and models is observed for the average event shape, while the event shape distributions exhibit large discrepancies. This is a very important message: the average measurements do not contain enough information, hence, care needs to be taken when extracting physics from models. For example, in the concrete case of the event generators reported in these references [24, 25], they overestimate significantly the contribution of back-to-back jets events and underestimate the contribution of isotropic events at high multiplicity.

3. Multi-parton interactions, event multiplicity and transverse spherocity

MPI are a theoretically [26] expected phenomenon, their effects have been observed by several experiments [27–29]. MPI in Pythia are crucial for the description of observables like multiplicity distributions, underlying event, correlations of the average transverse sphericity with multiplicity, and, together with color reconnection, to produce flow-like patterns.

Figure 1 shows the correlation between mid-rapidity charged hadron multiplicity and the number of multi-parton interactions ($N_{mpi}$) for $\sqrt{s} = 7$ TeV. Two features are observed. First, the width of the distribution for multi-parton interactions prevents using multiplicity alone as a selective parameter. Namely, in a given multiplicity interval there are events emanating from a very different number of MPI, hence of very different nature. The second feature is the tendency of saturation which is observed at high multiplicity. ALICE has measured the number of independent sources of particle production as a function of the event multiplicity using an approach based on two-particle azimuthal correlations, and it has reported a saturation effect at high multiplicity [30].
4. Results

4.1. Jet production as a function of multiplicity and transverse spherocity

As discussed earlier, event shapes are tools which allow the classification of the events according to their jet content. To probe the tool, this section shows the results of applying a jet finder to samples of events, where a pre-selection based on transverse spherocity was implemented. The jet finder Fast jet 3.0.6 [31] has been used, that implements the anti-$k_T$ algorithm with a jet radius of 0.4. The minimum jet $p_T$ was set to 10 GeV/$c$.

The left panel of Fig. 3 shows the average number of jets as a function of the event multiplicity. Results are presented for different $S_0$ intervals. As expected from transverse spherocity definition, for isotropic events ($S_0 > 0.8$) the average number of jets is below 1. This number increases when reducing transverse spherocity. Overall, the jet production rises with the event multiplicity, but the largest increase is seen in jetty-like events ($S_0 < 0.2$). A complementary study is presented in the right panel, where the average jet $p_T$ is plotted as a function of multiplicity. For isotropic events, the small fraction of jets which survives after the $S_0$ cut, exhibits an average $p_T$ which is flat and close to 10 GeV/$c$. On the contrary, a strong dependence is found for jetty-like events where at high multiplicity the average $p_T$ is above 30 GeV/$c$.

The same trend is observed if, instead of jets, inclusive charged hadrons are considered. In this case $\langle p_T \rangle$ in low transverse spherocity events shows a steeper rise with multiplicity than in the MB case. In contrast $\langle p_T \rangle$, for isotropic events shows a weaker multiplicity dependence. This suggests that the choice of narrow $S_0$ bins would allow studies with much less fluctuations.

4.2. Proton-to-pion ratio as a function of multiplicity and transverse spherocity

The particle ratios are among the important observables in nucleus interactions, for example, the proton yield normalized to that for pions encodes flow information and allows to test new hadronization mechanisms like quark recombination/coalescence [32]. Again, multiplicity alone is not able to display prominently the important features.
The left panel of Fig. 4 shows that the ratio as a function of $p_T$ is approximately the same for inelastic (minimum bias) and high multiplicity events. In fact, a weak flow-like effect is observed, namely, the depletion of the ratio at low $p_T$ for high multiplicity events with respect to the MB case. For comparison, in the right panel a selection using transverse spherocity is implemented. In fact, a weak flow-like effect is observed, namely, the depletion of the ratio at low $p_T$ for high multiplicity events. In this case the ratio for low $S_0$ events does not exhibit a “bump”. On the other hand, for isotropic events (large $N_{\text{pp}}$ implying stronger flow-like behavior) the ratio displays a similar behavior to that observed in LHC data [19]. At low $p_T$ there is a depletion when the multiplicity increases, then a crossing point is observed, in this case at $p_T \approx 2.5$ GeV/c. This crossing is followed by an enhancement and, at higher $p_T$ (> 8 GeV/c), the ratio returns to the value obtained for MB. This result carries two messages: the jetty-like events have, if any, a much lower flow-like effect than isotropic events; and the $p_T$ spectra and particle ratios should be extracted for both event classes in order to understand the role of jets when one studies, for example, radial flow.

4.3. Results from the blast-wave analysis and comparison with experimental data

It has been discussed in a previous letter [14] that the transverse momentum distributions of identified hadrons obtained from Pythia 8.180 exhibit flow-like features, the effect has been traced to color reconnection and multi-parton interactions. In heavy ion collisions the radial expansion can be extracted through a blast-wave analysis, where, a simultaneous fit to identified hadron $p_T$ distributions for each multiplicity bin is done. This parameterization assumes a locally thermalized medium, expanding collectively with a common velocity field and undergoing an instantaneous common freeze-out [33]. The simultaneous fit to all particle species under consideration can provide insight on the common kinetic freeze-out properties of the system. However, one needs to keep in mind that the values which come out from the fit depend substantially on the fit range.

In this section the results from the blast-wave analysis reported by ALICE for p-Pb collisions are compared with those obtained from a similar analysis applied to pp collisions simulated with Pythia 8.180 at $\sqrt{s} = 5.02$ TeV. The parameters extracted from the fits are studied as a function of event multiplicity and transverse spherocity. For the calculation of the transverse spherocity only primary charged hadrons with transverse momenta above 0.15 GeV/c and $|\eta| < 1$ are considered. The transverse momentum ($p_T$) spectra is (are) calculated counting primary charged hadrons within $|\eta| < 1$ ($|\eta| < 1$). The transverse momentum distributions of the particle species used in the blast-wave analysis are displayed in Table 1 along with their corresponding fit ranges. Figure 5 shows some examples of $p_T$ distributions at high multiplicity and for two extreme transverse spherocity event classes: jetty-like ($S_0 < 0.2$) and isotropic ($S_0 > 0.8$). Within 10% all the MC $p_T$ spectra are well described by the blast-wave model assuming a common transverse velocity $\langle \beta_T \rangle$ and temperature $T_{\text{kin}}$. Using somewhat different fit ranges for pions, kaons, protons and $\Lambda$; a similar observation is mentioned by the ALICE Collaboration in the analysis of p-Pb data [5].

Finally, Fig. 6 shows the evolution of $\langle \beta_T \rangle$ vs. $T_{\text{kin}}$ with event multiplicity and transverse spherocity. The results are compared with ALICE p-Pb data. The low transverse spherocity (jetty-like) events do exhibit a high $\langle \beta_T \rangle$ value.
Figure 5. (Color online). Transverse momentum distributions of identified hadrons for jetty-like (upper) and isotropic (bottom) events. Results for three different multiplicity classes are shown. A simultaneous blast-wave fit has been implemented, the results (lines) are plotted together with the \( p_T \) spectra.
with a low $T_{\text{kin}}$, while the high spherocity events exhibit a lower $\langle \beta_T \rangle$ and higher $T_{\text{kin}}$. This represents in retrospect an evidence: the jetty-like events by their nature will have a tendency to mimic flow, consequently, in radial flow studies care needs to be taken to avoid the non flow effects. The comparison of the simulations with the experiment demonstrates that in the treatment of the blast wave analysis one has to be careful to identify effects coming from jets from those from collective phenomena.

5. Conclusions

The transverse spherocity has been used to study the influence of the multiparton interactions on the final state in pp collisions. It was demonstrated that the number of multi-parton interactions is strongly correlated with the final event multiplicity at least up to a point where a saturation of the number of multi-parton interactions occurs. The transverse spherocity selection allows to identify and analyze two extreme cases: the jetty-like and the isotropic events situated at the two ends of the transverse spherocity spectrum. In the transverse spherocity event classes, narrower $N_{\text{mpi}}$ distributions are achieved than in the case without any selection on event shape. To illustrate the application of these variables, different studies were done, namely, jet production, identified particle ratios and blast-wave analysis. The results show the benefits of the combined multiplicity and event shape analysis. We conclude that a more widespread use of event shape variables in the analysis of the data may bring us a much better understanding of the detail of the collisions. The results using Pythia 8.180 simulations are qualitatively very similar to the experimental observations.

### Table 1. Fit ranges used in the blast-wave analysis.

| Particle species | Fit range |
|------------------|-----------|
| $\pi^+ + \pi^-$  | 0.5-1.0 GeV/c |
| $K^+ + K^-$      | 0.3-1.5 GeV/c |
| $K^0_S$          | 0.3-1.5 GeV/c |
| $p + \bar{p}$   | 0.8-2.0 GeV/c |
| $\phi$           | 0.8-2.0 GeV/c |
| $\Lambda$        | 1.0-2.1 GeV/c |
| $\Xi^- + \Xi^+$  | 1.2-2.6 GeV/c |
| $\Omega^- + \Omega^+$ | 1.3-2.8 GeV/c |
in p-Pb and Pb-Pb data (e.g., different particle composition in the jet and in the bulk regions). Suggesting that more differential studies using event shapes reveal interesting features which could be exploited to get physical information as well as to improve models used in the MC generators.

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References

[1] V. Khachatryan, et al., JHEP 1009 (2010) 091.
[2] S. Chatrchyan, et al., Phys. Lett. B718 (2013) 795–814.
[3] B. Abelev, et al., Phys. Lett. B719 (2013) 29–41.
[4] B. B. Abelev, et al., Phys. Lett. B726 (2013) 164–177.
[5] B. B. Abelev, et al., Phys. Lett. B728 (2014) 25–38.
[6] A. O. Velasquez, Nuclear Physics A932 (0) (2014) 146 – 151.
[7] P. Boz´ek, W. Broniowski, Phys. Rev. C88 (2013) 014903.
[8] P. Boz´ek, W. Broniowski, G. Torrieri, Phys. Rev. Lett. 111 (2013) 172303.
[9] A. Bzdak, G.-L. Ma, Phys. Rev. Lett. 113 (2014) 252301.
[10] G.-L. Ma, A. Bzdak, Phys. Lett. B739 (2014) 209–213.
[11] B. Schenke, S. Schlichting, R. Venugopal, Azimuthal anisotropies in p+Pb collisions from classical Yang-Mills dynamics. arXiv:1502.01331.
[12] A. Dumitru, K. Dusling, F. Gelis, J. Jalilian-Marian, T. Lappi, et al., Phys. Lett. B697 (2011) 21–25.
[13] T. Sjostrand, S. Mrenna, P. Z. Skands, Comput. Phys. Commun. 178 (2008) 852–867.
[14] A. Ortiz Velasquez, P. Christiansen, E. Cuautle Flores, I. Maldonado Cervantes, G. Pa´c, Phys. Rev. Lett. 111 (4) (2013) 042001.
[15] T. Sjostrand, Colour reconnection and its effects on precise measurements at the LHC. arXiv:1310.8073.
[16] F. Antinori, N. Armesto, P. Bartalini, R. Bellwied, P. Braun-Munzinger, et al., Thoughts on opportunities from high-energy nuclear collisions. arXiv:1409.2981.
[17] E. Cuautle, R. Jimenez, I. Maldonado, A. Ortiz, G. Pa´c, Disentangling the soft and hard components of the pp collisions using the spherocity approach. arXiv:1404.2372.
[18] B. Abelev, et al., Phys. Rev. C90 (2014) 054901.
[19] M. Veldhoen, Nucl. Phys. A910-911 (2013) 306–309.
[20] A. Zimmermann, Production of strange particles in charged jets in p-Pb and Pb-Pb collisions measured with ALICE at the LHC. arXiv:1502.01263.
[21] B. Abelev, et al., Phys. Rev. Lett. 109 (2012) 252301.
[22] A. Banfi, G. P. Salam, G. Zanderighi, JHEP 1006 (2010) 038.
[23] R. Corke, T. Sjostrand, JHEP 1103 (2011) 032.
[24] G. Aad, et al., Phys. Rev. D88 (3) (2013) 032004.
[25] B. Abelev, et al., Eur. Phys. J. C72 (2012) 2124.
[26] T. Sjostrand, M. van Zijl, Phys. Rev. D36 (1987) 2019–2041.
[27] S. Chekanov, et al., Nucl. Phys. B792 (2008) 1.
[28] F. Abe, et al., Phys. Rev. Lett. 79.
[29] V. Abazov, et al., Phys. Rev. D81.
[30] B. Abelev, et al., Journal of High Energy Physics 2013 (9).
[31] M. Cacciari, G. P. Salam, G. Soyez, Eur. Phys. J. C72 (2012) 1896.
[32] M. He, R. J. Fries, R. Rapp, Phys. Rev. C82 (2010) 034907.
[33] E. Schnedermann, J. Sollfrank, U. Heinz, Phys. Rev. C48 (1993) 2462–2475.