Abstract. I report results from the CLEO collaboration on semileptonic $B$ decays, highlighting measurements of the Cabibbo-Kobayashi-Maskawa matrix elements $|V_{cb}|$ and $|V_{ub}|$. I describe the techniques used to obtain the recent improvements in precision for these measurements, including the use of the $b \rightarrow s \gamma$ photon spectrum to reduce hadronic uncertainties in semileptonic $B$ decays.

The study of semileptonic $B$ meson decays allows measurement of the Cabibbo-Kobayashi-Maskawa (CKM) matrix elements $|V_{cb}|$ and $|V_{ub}|$, providing important inputs to a test of the unitarity of the CKM matrix, which governs the weak charged current and gives rise to $CP$ violation in the standard model. The rate for a $b$ hadron to decay weakly to hadrons containing a $c$ or $u$ quark is proportional to $|V_{cb}|^2$ or $|V_{ub}|^2$ respectively. The absence of final-state interactions in semileptonic decay make the interpretation less dependent on hadronic matrix elements than fully hadronic $B$ decays, although hadronic uncertainties still limit the precision of $|V_{ub}|$ and $|V_{cb}|$ measurements.

The current round of measurements from CLEO continues to test the hadronic calculations needed to disentangle weak matrix elements from strong interaction effects. For decays of $B$ mesons to exclusive final states, the hadronic effects are expressed in terms of a form factor that depends only on the momentum transfer $q^2$ to the lepton neutrino pair. By measuring decay rates as a function of $q^2$ we have begun to test the form factors, particularly for $b \rightarrow u \ell \bar{\nu}$ transitions. In decays to inclusive final states, under the assumption of parton-hadron duality, quark-level calculations may be compared to inclusive measurements to extract CKM matrix elements. Measurement of spectral distributions in inclusive decays gives additional observables to overconstrain theory parameters and test how well the theory and parton-hadron duality works.

$|V_{cb}|$ Measurements

CLEO has measured $|V_{cb}|$ using the decay $\bar{B} \rightarrow D^* \ell \bar{\nu}$ [1], where the decay rate as a function of $q^2$ is extrapolated to maximum $q^2$ where the $D^*$ is at rest in the frame of the initial $B$ meson. At this kinematic point the form factor $F$ is known to 4% of itself, owing to heavy quark symmetry considerations [2]. The differential decay rate is given

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by \( \frac{d\Gamma}{dq^2} = \frac{G_F^2}{48\pi^3}|V_{cb}|^2 \left[ \mathcal{F}(q^2) \right]^2 \mathcal{K}(q^2) \), where \( \mathcal{K} \) is a known kinematic function. Using \( \mathcal{F}(q_{\text{max}}^2) = 0.91 \pm 0.04 \) [3], we find \( |V_{cb}| = (47.4 \pm 1.4_{\text{stat}} \pm 2.0_{\text{syst}} \pm 2.1_{\text{fexpt}}) \times 10^{-3} \), somewhat higher than other results from \( \bar{B} \to D^* \ell \bar{\nu} \). The present world average from \( \bar{B} \to D^* \ell \bar{\nu} \) is \( |V_{cb}| = (42.4 \pm 1.2_{\text{exp}} \pm 1.9_{\text{theo}}) \times 10^{-3} \) [4].

A measurement of \( |V_{cb}| \) using the inclusive semileptonic decay rate is also possible. Here the experimental inputs are the branching fraction for \( \bar{B} \to X_c \ell \bar{\nu} \) and the \( B \) lifetime. The inclusive decay rate \( \Gamma_{\text{SL}} = \gamma |V_{cb}|^2 \), where \( \gamma \) comes from theory. Within the framework of heavy quark effective theory (HQET) [5], the inclusive semileptonic decay rate is expanded in a double series in \( \alpha_s^\prime \) and \( 1/M_n^\prime \), where \( M \) is the heavy quark mass. Hadronic effects enter both in the perturbative expansion and as expansion parameters, matrix elements of non-perturbative QCD operators. At \( O(1/M^2) \) there are two parameters: \( \lambda_1 \), which is proportional to the kinetic energy of the \( b \) quark in the \( B \) meson, and \( \lambda_2 \), which comes from the chromomagnetic operator. An additional parameter \( \Lambda \) relates the \( B \) meson mass to the \( b \) quark mass. From the \( B-B^* \) mass difference \( \lambda_2 = 0.128 \pm 0.010 \text{ GeV}^2 \). The other parameters can be estimated (e.g. in quark models) but they can also be measured using spectral moments in inclusive \( B \) decay. Moments, e.g. of the lepton energy spectrum, are also computed in HQET, allowing extraction of \( \lambda_1 \) and \( \Lambda \).

CLEO has a preliminary measurement of the inclusive semileptonic branching fraction using a high-momentum \( (p > 1.5 \text{ GeV/c}) \) lepton tag. The analysis is an update of Ref. [6]. The tag identifies a sample of \( B \) decays with high purity (98%). Additional electrons may come from the decay chain of the same \( B \) or from the decay of the other \( B \) meson in the event \( (e^+ e^- \to \Upsilon(4S) \to BB) \). Secondary leptons \( (b \to c \to \ell) \) and primary leptons are separated using kinematic and charge correlations, with a known correction from \( B^0\bar{B}^0 \) mixing. The new semileptonic branching fraction is \( 10.88 \pm 0.08 \pm 0.33\% \). The spectrum of electrons above 600 MeV is also obtained, from which spectral moments will be measured.

CLEO has recently measured spectral moments in inclusive semileptonic decay and in \( B \to X_s \gamma \). These are used to extract HQET parameters and reduce the theoretical uncertainty in inclusive \( |V_{cb}| \) measurements. CLEO measured the \( B \to X_s \gamma \) photon spectrum and moments (Fig. 1a) in [7]. In [8], CLEO measured the moments of the hadronic mass...
distribution in \( \bar{B} \to X_c \ell \bar{\nu} \) decays (Fig. 1b). Combining the constraints on \( \lambda_1 \) and \( \tilde{\Lambda} \) from the first moments of the photon energy and hadronic mass spectra, we obtain a solution for \( \lambda_1 \) and \( \tilde{\Lambda} \) and extract \( |V_{cb}| = (41.1 \pm 0.5_{\lambda_1, \tilde{\Lambda}} \pm 0.7_{\Gamma} \pm 0.8_{\text{HQET}}) \times 10^{-3} \) using the new CLEO branching fraction and PDG2003 lifetime average as inputs. The uncertainties from unknown \( \tilde{\sigma}(1/M^3) \) HQET parameters are dominant.

The lepton energy moments in \( \bar{B} \to X_c \ell \bar{\nu} \) are also sensitive to the HQET parameters, and CLEO has measured the lepton spectrum [9] and moments [10] above 1.5 GeV. From all of the moment measurements, one can assemble the constraints on the HQET parameters \( \lambda_1 \) and \( \tilde{\Lambda} \). Figure 1c shows the remarkable consistency of these measurements, lending credibility to the inclusive \( |V_{cb}| \) measurement.

At present the inclusive \( |V_{cb}| \) measurement is more precise (3%) than that from \( \bar{B} \to D^* \ell \bar{\nu} \), but with reliance on HQET for hadronic corrections. The first tests of HQET using spectral moments in inclusive \( B \) decays give us some confidence in the method, but additional tests with more inclusive moments are needed.

The agreement between inclusive and exclusive measurements is another test of our control of hadronic corrections. There is good agreement between inclusive and the world average exclusive \( |V_{cb}| \) measurements, but CLEO’s exclusive \( |V_{cb}| \) is larger than the inclusive measurement and other measurements using \( \bar{B} \to D^* \ell \bar{\nu} \).

### \( |V_{ub}| \) Measurements

Measurements of \( b \to u \ell \bar{\nu} \) have to contend with a 50–100 times larger background from \( b \to c \ell \bar{\nu} \). Requiring a lepton energy above the endpoint for \( b \to c \ell \bar{\nu} \) (\( \approx 2.3 \) GeV) is the easiest strategy to reduce background, but this cut near the edge of the spectrum introduces sensitivity to the motion of \( \bar{b} \) quark in the \( B \) meson. The sensitivity is reduced by using the \( b \to s \gamma \) photon spectrum [7], which is sensitive to the same hadronic effects at leading order [11, 12, 13, 14].

CLEO measured the lepton spectrum from \( B \) decays in the endpoint region \( E > 2.2 \) GeV and extracted a partial branching fraction of \( (2.30 \pm 0.15 \pm 0.35) \times 10^{-4} \) [15]. From the \( b \to s \gamma \) photon spectrum, the fraction of \( b \to u \ell \bar{\nu} \) events passing the lepton energy cut is \( f_u = 0.130 \pm 0.024 \pm 0.015 \). This gives \( |V_{ub}| = (4.08 \pm 0.34_{\text{exp}} \pm 0.44_{f_u} \pm 0.16_{\Gamma} \pm 0.24_{\text{NLO}}) \times 10^{-3} \), where the theoretical uncertainties are \( \Gamma \), from [16, 17], and NLO, from sub-leading terms relating hadronic effects in \( b \to u \ell \bar{\nu} \) and \( b \to s \gamma \).

CLEO has also measured \( |V_{ub}| \) in the exclusive modes \( \bar{B} \to |\pi/\rho/\omega/\eta| \ell \bar{\nu} \) [18], where kinematics from full reconstruction of the final state gives the needed suppression of \( b \to c \ell \bar{\nu} \). The neutrino is reconstructed from the missing energy and momentum of the event, taking advantage of CLEO’s large solid angle (95%). Combined with a lepton and light meson candidate, energy and momentum conservation leads to signal peaks in \( \Delta E = E - E_{\text{beam}} \) and \( M_{n\ell \bar{\nu}} \), with \( S/B \approx 1 \). We perform a simultaneous maximum likelihood fit in \( \Delta E \) and \( M_{n\ell \bar{\nu}} \) to seven sub-modes. Signals for \( \pi \) (Fig. 2b) and \( \rho \) are extracted separately in three \( q^2 \) bins. Given form factors from theory, we extract \( |V_{ub}| \) from a fit to \( d\Gamma/dq^2 \) (Fig. 2b). Combining \( \bar{B} \to \pi \ell \bar{\nu} \) and \( \bar{B} \to \rho \ell \bar{\nu} \) results we find \( |V_{ub}| = (3.17 \pm 0.17|_{\text{stat}} \pm 0.16|_{\text{syst}} \pm 0.53|_{\text{theo}} \pm 0.03|_{\text{FF}}) \times 10^{-3} \). This result uses form factors from Lattice QCD (\( q^2 > 16 \text{ GeV}^2 \)) and light cone sum rules (\( q^2 > 16 \text{ GeV}^2 \)) where each are...
most reliable. In a test of $\bar{B} \to \pi \ell \bar{\nu}$ form factors, ISGW2 \[19\] is disfavored (Fig. 2b).

We find good agreement between measurements of $|V_{ub}|$ using inclusive and exclusive techniques. The theoretical uncertainty on the form factor normalization currently limits the precision of the exclusive $|V_{ub}|$ measurement. In the future, unquenched Lattice QCD calculations can improve the $\bar{B} \to \pi \ell \bar{\nu}$ form factor in a limited region of $q^2$. The inclusive $b \to u \ell \bar{\nu}$ measurement can be further improved with increased $b \to s \gamma$ statistics and better phenomenological understanding of non-perturbative shape functions for the $B$ meson \[20\] \[21\] \[22\]. Comparison between inclusive measurements that use different kinematic cuts (more inclusive and away from the endpoint region) will increase our confidence in inclusive $|V_{ub}|$ measurements. Since the principal background comes from $b \to c \ell \bar{\nu}$, better knowledge of the dominant semileptonic $B$ decays will improve systematic errors for both inclusive and exclusive measurements.

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