ABSTRACT

By tracing the distribution of cool dust in the extended envelopes of post-asymptotic giant branch stars and (proto-)planetary nebulae ((P)PNs), we aim to recover, or constrain, the mass-loss history experienced by these stars in their recent past. The Far-Infrared Surveyor (FIS) instrument on board the AKARI satellite was used to obtain far-infrared maps for a selected sample of post-AGB stars and (P)PNs. We derived flux densities (aperture photometry) for 13 post-AGB stars and (P)PNs at four far-infrared wavelengths (65, 90, 140, and 160 μm). Radial (azimuthally averaged) profiles are used to investigate the presence of extended emission from cool dust. No (detached) extended emission is detected for any target in our sample at levels significant with respect to background and cirrus emission. Only IRAS 21046+4739 reveals tentative excess emission between 30′′ and 130′′. Estimates of the total dust and gas mass from the obtained maps indicate that the envelope masses of these stars should be large in order to be detected with the AKARI FIS. Imaging with higher sensitivity and higher spatial resolution is needed to detect and resolve, if present, any cool compact or extended emission associated with these evolved stars.

Key words: circumstellar matter – infrared: stars – planetary nebulae: general – stars: AGB and post-AGB

Online-only material: color figures

1. INTRODUCTION

During the asymptotic giant branch (AGB) phase, a low- and intermediate-mass star experiences mass loss (e.g., by thermal pulses) that acts as the main source of replenishment of gas and dust to the interstellar medium (ISM; Herwig 2005). When the AGB star evolves to the post-AGB phase (and later to the proto-planetary nebular (PPN) and the PN phase) the mass loss drops by several orders of magnitude, while the circumstellar shells continue to drift away (pushed by the—episodic—(post)-AGB wind) from the star into the surrounding medium (Van Winckel 2003). For a constant mass-loss rate a smooth envelope with decreasing column density, inverse to the radial distance from the central star, is expected. If, on the other hand, the mass-loss rate fluctuates (as the star goes through thermal pulses during the AGB; Vassiliadis & Wood 1993), then enhanced emission should be observed in the form of multiple dust shells around the central star. Thus, the history of mass loss during the AGB phase is imprinted on the dust shell of the PN until it is disrupted by the fast post-AGB wind. Pre/proto-planetary nebulae (PPNe) are objects in transition from the high mass-loss rate AGB phase to that of an ionized PN. However, little is known yet about the large extended shells around these evolved stars. Therefore, young PNe and proto-PNe may have outer shells in which the fossil record of the mass loss during the AGB phase is still preserved. Cool dust present in circumstellar shells emits mainly in the infrared. Early studies suggested that many PNe are extended in the infrared, but usually with sizes comparable to the optical diameter (Hawkins & Zuckerman 1991). This indicated that the dust grains are generally well mixed with the ionized gas. However, IRAS also showed a few PNe with the far-infrared emission extending beyond the limits derived from optical images. This suggested that cool dust may be present in a neutral envelope outside the inner ionized zone. The spatial extent and relative brightness distribution of the mid- and far-infrared radiation of the dust shells around PNe provide important constraints on the evolution of the central star. In particular, the mass-loss conditions in the AGB strongly constrain evolutionary models as this affects the distribution of ionized gas, molecular gas, and cool dust in the resulting PN.

Recent Spitzer/MIPS observations by Do et al. (2007) have not revealed circumstellar dust shells around CRL 2688, as reported with Infrared Space Observatory (ISO). The latter were reported by Speck et al. (2000) who identified dust shells of about 2–3 pc in two post-AGB stars. Klaas et al. (2006) tentatively inferred a slightly extended emission from PNe NGC 7008 by means of ISOPHOT multi-aperture photometry. Spitzer/MIPS observations of PN NGC 2346 and NGC 650 (at 24, 70, and 160 μm; Su et al. 2004 and Ueta 2006, respectively) both show the 24 μm morphology is strikingly similar to that seen in the optical (due to the contribution of nebular emission lines in this band), while the 70 and 160 μm images show an extended (~90′′) cold dust envelope around the nebula. More recently, Chu et al. (2009) compared the 24 μm surface brightness profiles with those in the Hα line for a sample of 36 Galactic PNe; they found that the infrared emission can be more extended, similar or present in the shell center only. Chu et al. (2009) suggested that these different groups of PNe have an evolutionary connection—with the youngest PNe showing more extended 24 μm emission (due to hot dust) than the Hα emission. First results with Herschel reveal detached and extended emission in the far-infrared associated to evolved stars (Groenewegen et al. 2011) as well as wind–ISM interaction at several arcminutes from some objects, such as CW Leo (Ladjal et al. 2010). The latter was also traced in the FUV by Galaxy Evolution Explorer (Sahai et al. 2010) similar to the shock...
interface observed for Mira (Martin et al. 2007). Kerschbaum et al. (2010) reported detached dust shells (at ∼30′′–50′′) associated with three AGB stars located at 260–825 pc (we note that the most distant AGB star had the dust shell located farthest away from the point source). Extended—parsec-sized—far-infrared emission is expected to be present outside the ionized region in the neutral zone around post-AGB stars and PNe as indicated by observations (Speck et al. 2000), as well as theoretical work (Villaver et al. 2002). Therefore, if indeed the typical sizes of far-infrared dust emission extends over distances of 2–3 pc, the apparent sizes would range from 1′ to 10′ for distances in the range 0.5–3 kpc. These are the angular scales that the most distant AGB star had the dust shell located farthest away from the point source. Extended—parsec-sized—far-infrared emission will also depend strongly on the transitory correction and applied a smooth filter (90s boxcar) for the long-term changes in detector responsivity, cosmic-ray glitches, and ionizing radiation spikes. The distortion of the FIS array detectors’ field of view and their alignment is also corrected. The image background level was set to zero by subtracting the sky background in the time domain. The focal-plane coordinate system was converted into the ecliptic coordinate system. In this way, all frames have the same orientation since the scanning is done parallel to the ecliptic plane. The images are created with the AKARI FIS images. The exact size and shape of the extended emission will also depend strongly not only on the mass-loss rate but also on the relative velocity of the star and the density of the surrounding ISM.

In this paper, we present a far-infrared survey of 13 post-AGB stars and (young/proto) PNe obtained with the Far-Infrared Surveyor (FIS) on the Japanese Infrared Astronomical Mission AKARI. The properties of the far-infrared emission of these objects and their immediate surroundings are reported and discussed. In Section 2, we discuss the observations and the data processing. The main observational results, including photometry, are reported in Section 3. In Section 4, we discuss the main results of this work and the limits on the possible presence of extended cold dust emission associated with the post-AGB objects in our sample. Section 5 concludes the paper.

2. INFRARED MAPPING WITH THE FAR-INFRARED SURVEYOR

We used the FIS instrument on board the AKARI infrared astronomy mission (Murakami et al. 2007) to obtain pointed observations of 13 post-AGB stars and (proto) planetary nebulae (PNe). The FIS (Kawada et al. 2004, 2007) provides scan maps at four wavelength filter bands with the short wavelength (65 and 90 μm) and long wavelength (140 and 160 μm) detectors. The FIS was used in the FISO1 compact source photometry scan-mode. In this mode, the detectors sweep the sky twice in round trips for each pointing. The scan speed was set to 15′′ s−1 and each scan takes 157.5 s to complete. Between two round trips, the scan path is shifted by either 70′′ or 240′′ resulting in final maps of 10′ × 40′ or 13′ × 40′, respectively. The sampling (detector readout) mode was either nominal (with reset interval 0.5 or 1.0 s) or in CDS mode (which takes the ramp differential data: i.e., the difference between the begin and end voltage for each interval). The settings for each pointed observation are given in the observational log (Table 1) and basic target information (ID/names, type, FIS coordinates) given in Table 1.

The pipeline modules provided by the AKARI team (AKARI FIS Data User Manual Version 1.2; Verdugo et al. 2007) were used for the processing of the FIS data. This includes dark subtraction, responsivity correction, flat fielding, and correction for slow transients (e.g., after calibration lamp exposures). The pipeline also takes care of removing gradual long-term changes in detector responsivity, cosmic-ray glitches, and ionizing radiation spikes. The distortion of the FIS array detectors’ field of view and their alignment is also corrected. The image background level was set to zero by subtracting the sky background in the time domain. The focal-plane coordinate system was converted into the ecliptic coordinate system. In this way, all frames have the same orientation since the scanning is done parallel to the ecliptic plane. The images are created with grid sizes of 30′ for the Wide-L and N160 bands and 15′ for the Wide-S and N60 bands. We verified that the positions of the observed targets (obtained via a two-dimensional Gaussian fit to the infrared point source) agree within ∼10′′ with those of the optical/near-infrared counterparts (e.g., Two Micron All Sky Survey/Simbad: Skrutskie et al. 2006; Kerber et al. 2003). The nominal FWHM of the point-spread function is ∼40′′ for the LW and ∼30′′ for the SW detectors, respectively. Relevant filter and detector properties are given in Table 2. More details on the FIS and its performance are given in Kawada et al. (2004, 2007) and Shirahata et al. (2009).

The nominal FWHM of the point-spread function is ∼40′′ for the LW and ∼30′′ for the SW detectors, respectively. Relevant filter and detector properties are given in Table 2. More details on the FIS and its performance are given in Kawada et al. (2004, 2007) and Shirahata et al. (2009). In effect, for the LW we applied the transitory correction and applied a smooth filter (90s boxcar) with 1.5σ-clipping. For the SW we applied in addition (for the non-CDS observations) a median filter (200s) and used the local flat. The reduced images for all targets are shown in Figure 1.
Figure 1. The final images in the 65, 90, 140 and 160 μm band for each target are shown in subfigures (a) to (m). Contour levels are drawn in logarithmic intervals (scaled between minimum and maximum brightness). The crosses show the source central position. All images are oriented with the scan direction upwards (i.e., along the ecliptic) which facilitates comparisons. Pixel sizes are 15′′ and 30′′ for the SW and LW bands, respectively (Note that the axes labels are in pixel units). Also shown are the aperture and sky annuli used for the aperture flux photometry. Note that the strong bar (angle of ∼50° with respect to the ecliptic) feature in the SW band is a detector artifact (see e.g., Shirahata et al. 2009). It is not present in the LW images. The subfigures show (a) IRAS 07134+1005; (b) IRAS 09425-6040; The 140 and 160 μm images reveal a cold dust emission feature, with peak temperature around 20 K, located 10′ upwards (in the direction parallel to the ecliptic) of the central source. Potentially, this is an old detached shell, but we cannot rule out diffuse cirrus emission; (c) IRAS 10178-5958; (d) IRAS 10197-5750; (e) IRAS 11478-5654; (f) IRAS 13428-6232; (g) IRAS 15318-7144; (h) IRAS 17119-5926; (j) IRAS 17395-0841; (k) IRAS 19500-1709; (l) IRAS 20104+1950; (m) IRAS 22036+5306. (A color version of this figure is available in the online journal.)

3. RESULTS AND DATA ANALYSIS

From the reduced images, we extract information on the total point-source flux density and detection limits on the extended emission. Azimuthally averaged radial profiles were constructed in order to search for extended spherical structures associated with faint extended dust emission.

3.1. Absolute Point-source Aperture Photometry and Detection Limits on the Extended Emission

Circular aperture photometry was applied to all sources to derive absolute fluxes. For the SW fluxes, the aperture radius was set to 5 pixels (or ∼75″). The sky background was determined from an annulus with inner and outer radii of 9 and
12 pixels, respectively. For the LW flux measurements, we adopted an aperture radius of 3 pixels (∼90″). The sky background was computed from an annulus with inner and outer radii of 6 and 8 pixels, respectively. To convert the obtained image fluxes to absolute fluxes, we apply the diffuse-to-point correction factor as well as the aperture correction factor (see Table 2 and Shirahata et al. 2009). Note that the flux calibration by aperture photometry refers to a flat (νFν = constant) spectrum while the real flux depends on the spectral energy distribution (SED) of the source. At this stage, we have not applied any color correction. Shirahata et al. (2009) report an absolute FIS calibration accuracy better than 15%, except for the longest wavelength (∼50%). These authors also report a tentative 10% decrease in observed flux in all bands for the fast (15″ s⁻¹) scan speed with respect to the nominal 8″ s⁻¹ scan speed. As all our targets were obtained with 15″ s⁻¹, this may result in a small underestimation of the final flux densities. Furthermore, we assume these calibration factors are valid also for CDS mode (recommended for very bright sources) observations although this has not been verified in depth. Our results can be compared directly with the SED provided for five targets included in the Toruń catalogue of Galactic post-AGB and related objects (Szczepanek et al. 2007). From these SEDs, there appears to be no evidence for systematic lower (or
Figure 1. (Continued)

higher) than expected flux densities. Figure 2 illustrates that the absolute fluxes for bright point sources measured with AKARI FIS at 65 and 90 $\mu$m are, at a first approximation, consistent with the IRAS fluxes at 60 and 100 $\mu$m, even without applying a color correction. For the post-AGB/proto-PNe in our sample, nebular lines are expected to contribute little to the FIR broadband emission; this effect could be more significant for the PNe. This consistency is also an important calibration check for flux densities reported in the AKARI bright source catalog (Yamamura et al. 2010).

For the 90 and 140 $\mu$m maps, we also measure the average flux noise level (Jy arcmin$^{-2}$) from the sky background within 2′–5′ from the point source. The 3$\sigma$ detection limits derived from our FIS maps for each target are listed in Table 3. These upper limits are used in Section 4 to derive upper limits on the dust and gas mass in the extended envelopes.

3.2. Azimuthally Averaged Radial Intensity Profiles

We extracted the azimuthally averaged radial profile for each target in each filter band. In addition, we extract radial profiles from calibration point sources (Vesta and Neptune) processed in the same manner as our targets. The radial profiles for the calibration point sources and all post-AGB and PNe sources are shown on a log–log scale in Figure 3 for the two most sensitive bands at 90 and 140 $\mu$m. Due to the width of the image maps, proper radial profiles can only be obtained out to radii of about 300″. As indicated above, there is no unambiguous
evidence for extended emission within 2′ of the point source. The sensitivity of these maps is insufficient to unambiguously detect weak detached emission between 2′ and 5′. One potential exception is IRAS 21046+4739, for which localized far-infrared emission at 140 μm (and also present in the 160 μm map) could be related directly to the PN.

4. DISCUSSION

Individual modeling of the SED of these targets with a complex radiative transfer code (like DUSTY) is beyond the scope of this paper, which is focused on the cold extended/detached circumstellar dust emission, as the new photometry points (of the central source) do not add significant constraints to the problem with respect to previous studies. Furthermore, the comparison of our AKARI FIS maps with simulated images (or SEDs) of the expected extended dust emission seems to be inappropriate due to the lack of a clear detection of extended emission in the AKARI images; such models (e.g., 2-DUST; Ueta & Meixner 2003) also have their own limitations and assumptions, as well as many free parameters. In particular, the uncertain distances to the post-AGB objects in our sample affect the model results. Thus, we adopt here a simple dust model, together with the upper limits derived from the new
FIS maps (Table 3), in order to give first-order estimates on upper limits on the detected dust (and gas) mass in the extended envelopes of these evolved objects.

4.1. Upper Limits on the Detected Dust+Gas Mass in Extended Envelopes

Interstellar and circumstellar dust emits primarily in the infrared. Assuming that the observed flux comes from an optically thin emission region, the total mass of the dust and the associated gas can be estimated, with a few assumptions about the basic properties of the dust, from the following equation (see, e.g., Li 2005; Ueta et al. 2010):

$$M_{\text{gas+dust}} [g] \approx \frac{d^2 F_\nu \lambda^2}{2 k T_{\text{dust}} \kappa_{\text{abs}}(\lambda)} g 2d,$$

(1)

with the flux $F_\nu$ in erg cm$^{-2}$ s$^{-1}$ Hz$^{-1}$ ($F_\nu = 10^{-23} F(\text{Jy})$), the distance $d$ in cm, the wavelength $\lambda$ in cm, the dust temperature $T_{\text{dust}}$ in K, the specific opacity or dust mass absorption coefficient, i.e., opacity, $\kappa_{\text{abs}}(\lambda)$ in cm$^2$ g$^{-1}$, and the Boltzmann constant in erg K$^{-1}$. For the gas-to-dust mass ratio, $g 2d$, we take 160 for oxygen-rich stars and 400 for carbon-rich
stars (see, e.g., Knapp 1985; Heras & Hony 2005). Apart from the uncertainties in the distance, large uncertainties result from poor knowledge of grain properties as expressed through $\kappa_{\text{abs}}$ (Li 2005). $F_\nu$ is the 3$\sigma$-upper limit on the flux, at 90 or 140 $\mu$m, detectable in a shell of 3 pc diameter around the central star (Villaver et al. 2002). We assume a homogenous distribution across an annulus with an inner radius of $1\arcmin$ and outer radius of $(5157/d)'$, respectively. The flux noise levels measured from the maps are given in Table 3. For the distance we take values either from the literature or, where this is unavailable, we assume $d = 1$–2 kpc (Table 4).

Values for opacity $\kappa$ are strongly dependent on (the composition of) the dust model or laboratory analog and can differ by an order of magnitude between studies. For C-rich post-AGB stars with polycyclic aromatic hydrocarbon (PAH) features that have a high degree of hydrogenation a common choice is either BE-type hydrogenated amorphous carbon dust from Colangeli et al. (1995) or the silicate–graphite–PAH dust model from Li & Draine (2001). The former gives $\kappa(90 \mu m) = 976$ cm$^2$ g$^{-1}$, while the latter gives, via $\kappa_{\text{abs}} = 2.92 \times 10^3 \cdot \lambda^{-2}$, $\kappa(90 \mu m) = 36$ cm$^2$ g$^{-1}$. Conservative dust mass limits are obtained by adopting the latter values for $\kappa$. These are also taken if we do not know...
the chemistry of the envelope. For the sources with an oxygen-rich dust or mixed chemistry in the envelope, we adopt $\kappa_{\text{abs}}$ for astronomical silicates (i.e., 63 cm$^2$ g$^{-1}$ at 90 $\mu$m; Draine & Lee 1984).

In computing the total extended dust and gas mass, we consider relatively cold dust grains with $T_{\text{dust}} = 20$ K as these correspond to old mass-loss events which has already cooled down and drifted farthest away (e.g., García-Hernández et al. 2006). The resulting upper limits to the “hidden” dust and gas mass in the extended shells around the observed stars are reported in Table 4. From Equation (1), we remark that the derived upper limit on the envelope mass reduces proportionally to an increase in dust grain temperature and/or the dust opacity. The dust mass limits are proportional to the adapted gas-to-dust ratio and the distance squared. In particular, the latter can introduce significant uncertainties.

The 90 $\mu$m FIS scan maps together with the aforementioned adopted grain properties provide upper limits to the total dust and gas mass between 0.1 and 7.5 $M_\odot$ (excluding IRAS 10197-5750 and IRAS 13428-6232 which show high levels of background noise). Noting the limits of the dust model (see above), we find that the dust and gas mass limits compare well with the total mass ($\sim$1 to 6 $M_\odot$) predicted to be lost by a star, with a certain initial mass, at the end of the AGB phase, as predicted by Villaver et al. (2002). The estimated dust and gas mass limits are generally high, indicating that the envelope masses of these evolved stars should have been large in order to have been clearly detected with the FIS. Nonetheless, for several cases the extended emission can be constrained to less than 1 $M_\odot$.

4.2. Emission Features in the FIS Maps

We briefly highlight several tantalizing, although tentative, features that can be discerned in the AKARI/FIS scan maps of several individual objects. The LW maps of IRAS 09425-6040 show an arc-like dust emission structure 10$^\prime$ upward (in the direction of the ecliptic) of the central source. It is also visible, though very faint, in the SW image. This structure is brightest at 140 $\mu$m, indicating a dust temperature of $\sim$20 K. The 140 and
### Table 2

| Band       | N| 60 | Wide-S | Wide-L | N| 160 |
|------------|---|----|--------|--------|---|----|
| \( \lambda \) (\( \mu m \)) | 65 | 90 | 140 | 160 |
| \( \Delta \lambda \) (\( \mu m \)) | 22 | 38 | 52 | 34 |
| Wavelength range (\( \mu m \)) | 50–80 | 60–110 | 110–180 | 140–180 |
| FWHM (\( ^\circ \)) | 32 \pm 1 | 30 \pm 1 | 41 \pm 1 | 38 \pm 1 |
| Pixel size (\( ^\circ \)) | 26.8 | 26.8 | 44.2 | 44.2 |
| Diffuse-to-point | 0.698 (flux)\(^{−0.0659} \) | 0.700 (flux)\(^{−0.0757} \) | 0.560 | 0.277 |
| Aperture correction | 0.882 | 0.865 | 0.881 | 0.910 |

**Notes.** See Verduzo et al. (2007) for additional details. Also listed are the adopted correction factors to obtain absolute flux levels. Aperture correction values are given for aperture radii of \( \sim 75'' \) and \( \sim 90'' \) for the SW and LW, respectively (from Shirahata et al. 2009).

### Table 3

| IRAS Name | \( F(65 \mu m) \) | \( \sigma_F \) | \( F(90 \mu m) \) | \( \sigma_F \) | \( F(140 \mu m) \) | \( \sigma_F \) | \( F(160 \mu m) \) | \( \sigma_F \) | 3\( \sigma \) Detection Limit |
|-----------|-----------------|-------------|-----------------|-------------|-----------------|-------------|-----------------|-------------|-------------------|
| 07134+1005 | 54.0 | 1.3 | 29.1 | 0.3 | 5.9 | 0.5 | 5.0 | 1.5 | 0.4 | 0.3 |
| 09425–6040 | 12.0 | 0.9 | 6.6 | 0.2 | 2.1 | 0.4 | \ldots | \ldots | 0.4 | 0.4 |
| 10178–5958 | 66.4 | 1.4 | 51.9 | 0.6 | (18.5) | 2.7 | (13.9) | 5.5 | 0.7 | 1.3 |
| 10197–5750 | 535.6 | 2.5 | 308.4 | 2.7 | (48.6) | 9.3 | \ldots | \ldots | 5.1 | 14.1 |
| 11478–5654 | 36.1 | 0.1 | 31.9 | 0.1 | 6.7 | 0.4 | 5.6 | 0.9 | 0.2 | 0.2 |
| 13428–6232 | 410.1 | 3.1 | 227.3 | 2.4 | \ldots | \ldots | \ldots | \ldots | 5.4 | 17.6 |
| 15318–7144 | 55.2 | 1.4 | 34.1 | 0.3 | 10.2 | 0.4 | 7.9 | 2.7 | 0.6 | 0.3 |
| 17119–5926 | 2.9 | 0.1 | 1.9 | 0.1 | 0.8 | 0.2 | \ldots | \ldots | 0.06 | 0.2 |
| 17395–0841 | 5.9 | 0.1 | 4.6 | 0.1 | (1.3) | 0.5 | \ldots | \ldots | 0.1 | 0.4 |
| 19500–1709 | 49.9 | 0.9 | 32.3 | 0.2 | 7.3 | 0.2 | 4.0 | 0.6 | 0.5 | 0.3 |
| 20104+1950 | 7.7 | 0.1 | 6.1 | 0.1 | 2.8 | 0.3 | (2.5) | 0.5 | 0.08 | 0.3 |
| 21046+4739 | 36.7 | 0.1 | 33.9 | 0.1 | 15.5 | 0.6 | 12.1 | 0.9 | 0.2 | 0.8 |
| 22036+5306 | 113.8 | 0.2 | 82.7 | 0.1 | 24.9 | 0.3 | 19.2 | 0.7 | 0.2 | 0.3 |

**Notes.** Values in parenthesis are less secure due to possible low-level “cirrus” contamination. The 3\( \sigma \) detection limits on the extended flux are in Jy arcmin\(^−2 \).

### Table 4

| IRAS Name | \( d \) (kpc) | Chemistry/Dust | Dust+Gas Mass Upper Limit (\( M_\odot \)) |
|-----------|---------------|-----------------|----------------------------------|
| 07134+1005 | 2.4\(^{(1)} \) | C | 1.1 |
| 09425–6040 | 1–2 | O | 0.2–0.88 |
| 10178–5958 | 1–2 | \ldots | 1.9–7.5 |
| 10197–5750 | 1–2 | O | 3.0–12.0 |
| 11478–5654 | 0.6\(^{(2)} \) | C | 0.39 |
| 13428–6232 | 1–2 | \ldots | 13.9–55.7 |
| 15318–7144 | 0.9\(^{(2)} \) | O | 0.31 |
| 17119–5926 | 1–2 | \ldots | 0.16–0.64 |
| 17395–0841 | 1.0\(^{(3)} \) | O | 0.08 |
| 19500–1709 | 1–2 | C | 1.3–5.4 |
| 20104+1950 | 1.8\(^{(2)} \) | C | 0.21 |
| 21046+4739 | 1.3\(^{(2)} \) | C | 0.53 |
| 22036+5306 | 2\(^{(4)} \) | O | 0.12 |

**Notes.** Flux limits are obtained from the observed FIS detection limits and the angular size of the envelope assuming a radius of 3.0 pc (approximately between 5' and 10') and excluding the inner 1'. See the text for details on the calculations and the adopted dust properties. In particular, we note the dependence of the dust mass with the distance squared. We adopt gas-to-dust ratios of 160 for oxygen-rich dust and 400 for carbon-rich dust (see the text). To obtain mass limits for a more generic gas-to-dust ratio of 200, divide the carbon-dust limits by 2.

**References.** (1) Hony et al. 2003; (2) Lenzunzi et al. 1989; (3) Gauja et al. 2003; (4) Sahai et al. 2008.

160 \( \mu m \) maps IRAS 10178–5958 show some extended emission in two arcs 4' downward and 8' upward of the central source. There is some indication of this emission at 90 \( \mu m \) but not at 65 \( \mu m \), indicating that this is cooler dust, with a temperature of less than 20 K. The 140 \( \mu m \) flux toward IRAS 10197–5750 may be underestimated due to extended “cirrus” emission contaminating the background measurements. The structure in the two SW bands is very similar (two arcs at 10' distance from the central source) while very different from the two LW bands in which a dust bar perpendicular to the direction of the ecliptic is observed. Thus, there appear to be two distinct structures that exhibit different average dust temperatures. The FIS maps of IRAS 13428–6232 show a wealth of structure in the far-infrared emission that show a higher flux at longer wavelengths indicating a lower temperature of this diffuse “cirrus” dust. The central source is only detected in the SW bands. The LW bands show an indication of a central point source but the extended structure makes it all but impossible to do photometry. Whether these structures are detached shells at about 7' and excluding the inner 1'.
far-infrared emission excess in the LW radial profile at radii between 30" and 130".

4.3. Discrepant Far-infrared Point-source Flux Densities

The FIR flux densities of IRAS 17119-5926 and IRAS 17395-0841 are significantly lower (by a factor two) than predicted from the IRAS data. The FIS images show no obvious source of contamination in the SW bands just as IRAS does not give a high value for cirrus emission. IRAS 17119-5926, however, is a peculiar source that is thought to have evolved only recently, in the last few decades, from a post-AGB into a PN (Parthasarathy et al. 1993). Therefore, variability in its infrared flux might occur due to recent mass-loss event(s) or evolution of the central star. This is supported also by the lower continuum level from the short wavelength spectrometer spectrum. The latter appearing in fact more consistent with the new FIS flux densities. For IRAS 17395-0841, the emission in the LW images does not appear to be directly related to the point source. The IRAS Point Source Catalog (PSC) lists a relatively high value for the cirrus emission at 100 μm. Four other sources in our sample have high cirrus values in the IRAS PSC but their FIS fluxes match well with those given by IRAS. Also, from the SEDs we find that nebular emission, as expected for post-AGB/proto-PNe, contributes very little to the observed broadband flux.

5. CONCLUDING REMARKS

Our observations provide imaging of post-AGB and (P)PNe beyond 100 μm for which only limited information is currently available. The AKARI FIS photometry is generally consistent, considering also color correction, differences in filter curves, and absolute calibration uncertainties, with the available SEDs, based in the infrared—primarily on IRAS.
and ISO data. We also conclude that the far-infrared point-source emission detected toward our sample of post-AGB stars and (P)PNe originates from a compact region unresolved by the FIS.

The observed background emission (measured for each map) was used to derive first-order upper limits on the mass in “cold” extended emission. The derived dust and gas masses are strongly dependent on, in particular, the adopted distance, and adopted grain size and emissivity (both of which affect the equilibrium extended emission near the central object at 140 and 160 \(\mu m\)). IRAS 10178-5958 shows two arcs at 4’ and 8’ from the central object at 140 and 160 \(\mu m\), IRAS 10197-5750 shows two shells at 10’ in the short wavelength filters and a peculiar cold dust bar at the longest wavelengths, and IRAS 21046+4739 reveals both in the FIS images and the radial profiles extended emission near the central source (<3’) at 140 and 160 \(\mu m\).

Further studies of the extended emission (in particular for the objects with tentative detections of mass-loss events) will require multi-wavelength observations to obtain not only an estimate of the total dust and gas mass but also some insight into the basic properties (temperature, size, composition) of grains in these extended shells. Imaging with both improved sensitivity and higher spatial resolution is required to observe extended structures due to cold dust associated with the sources presented above. In particular, the capabilities of the Herschel Space Observatory are favorable to detecting the coldest dust components in extended areas and at the high spatial resolution required to disentangle circumstellar dust from cirrus dust.

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