Fermiology of a topological line-nodal compound CaSb$_2$ and its implication to superconductivity: Angle-resolved photoemission study

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We performed angle-resolved photoemission spectroscopy with microfocused beam on a topological line-nodal compound CaSb$_2$ which undergoes a superconducting transition at the onset $T_c \sim 1.8$ K, to clarify the Fermi-surface topology relevant to the occurrence of superconductivity. We found that a three-dimensional hole pocket at the $\Gamma$ point is commonly seen for two types of single-crystalline samples fabricated by different growth conditions. On the other hand, the carrier-doping level estimated from the position of the chemical potential was found to be sensitive to the sample fabrication condition. The cylindrical electron pocket at the $\gamma(C)$ point predicted by the calculations is absent in one of the two samples, despite the fact that both samples commonly show superconductivity with similar $T_c$'s. This suggests a key role of the three-dimensional hole pocket to the occurrence of superconductivity, and further points to an intriguing possibility to control the topological nature of superconductivity by carrier tuning in CaSb$_2$.

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I. INTRODUCTION

Initiated by the discovery of topological insulators (TIs) whose surface or edge hosts gapless states despite the insulating nature of bulk [1–3], the search for new quantum states of matter characterized by the nontrivial topology is becoming one of major challenges in condensed-matter physics. While TI and topological superconductor (TSC)—a superconducting analog of TI—commonly contain a finite energy gap in the bulk (band gap and superconducting gap, respectively), the classification of band topology has been successfully extended to the systems with gapless bulk excitations, leading to the discovery of topological semimetals (TSMs). The first concrete example of TSMs is the three-dimensional (3D) topological Dirac semimetal such as Na$_3$Bi and Cd$_3$As$_2$ [4–8] where a linearly dispersive bulk valence band (VB) and a conduction band (CB) contact each other at a discrete point in the 3D $k$ space. Weyl semimetal is also characterized by the point touching of VB and CB, while the bulk band degeneracy is lifted by the breaking of time-reversal or space inversion symmetry [9–13]. There is another type of TSM called line-nodal semimetal where the crossing point extends one-dimensionally in the $k$ space (nodal line or loop), being protected by the crystal symmetry such as mirror reflection or nonsymmetric glide mirror symmetry [14–28]. While such TSMs, or more widely, materials containing Dirac/Weyl or line nodes in the bulk-band structure, are already a useful platform to realize exotic quantum phenomena such as unusual magnetotransport properties, the coupling of nodal electrons to superconductivity is another intriguing target of intensive investigations because such a coupling would provoke the topological superconductivity [29]. However, this essential issue remains highly unexplored, because superconducting materials whose low-energy excitations are characterized by Dirac/Weyl or line-nodal electrons are generally rare [16,30–32].

Recently, the superconductivity with superconducting transition temperature ($T_c$) of 1.7 K was discovered in a line-nodal...
In this paper, we report microfocused angle-resolved photoemission spectroscopy (μ-ARPES) of CaSb₂ to clarify the 3D band structure and Fermi-surface (FS) topology. We observed two different types of FSs in CaSb₂; one is a 3D hole pocket at the Γ point and another is a cylindrical electron pocket at the Y/Z points which are related to the 3D pocket. Such band crossings are connected to the line nodes at the k₀ = π plane and one of them is associated with the shallow electron pocket crossing Eᶠ, likely contributing to the transport properties, while these line nodes and Dirac nodes had to await experimental verification.

Magnetotransport measurements of CaSb₂ clarified a resistivity plateau under a weak magnetic field at 10 K and nonsaturating giant magnetoresistance, consistent with the existence of line nodes [34]. Electrical resistivity measurements under pressure revealed the dome-shaped Tc variation without Lifshitz transition in the normal state [35] and specific heat measurements signified the deviation from the BCS behavior [36], suggestive of the unconventional nature of superconductivity. On the other hand, ¹²¹⁷Sb nuclear quadrupole resonance (NQR) measurements suggested an exponential decrease in the inverse relaxation rate 1/T₁ at low temperature, supportive of the conventional s-wave superconductivity with a full gap [37]. The presence of cylindrical electron pockets with line nodes has been discussed in favor of topological superconductivity [33] associated with the dominant interorbital pairing interaction with odd-parity pairing within the cylindrical FSs, similar to the case of doped TIs [38] and doped Dirac semimetals [29,39]. Topological superconductivity was also predicted based on the full computation of the Z₂-enriched symmetry indicators [40]. Despite such intensive experimental and theoretical studies, there are no experimental reports on the band structure of CaSb₂. It is thus of great importance to experimentally clarify the electronic states of CaSb₂ to establish the interplay among nodal electrons, superconductivity, and topology.

In this work, we report microfocused angle-resolved photoemission spectroscopy (μ-ARPES) of CaSb₂ to clarify the 3D band structure and Fermi-surface (FS) topology. We observed two different types of FSs in CaSb₂: one is a 3D hole pocket at the Γ point and another is a cylindrical electron pocket at the Y/Z points. Such band crossings are connected to the line nodes at the k₀ = π plane and one of them is associated with the shallow electron pocket crossing Eᶠ, likely contributing to the transport properties, while these line nodes and Dirac nodes had to await experimental verification.

We prepared two different types of single crystalline samples for ARPES measurements, called here sample A and
sample B. The samples A and B were grown by a self-flux method with molar ratios of Ca:Sb=1:5 and 1:3.1, respectively. For sample A, Ca (Sigma-Aldrich, 99.9999%) and Sb (Alfa Aesar, 99.9999%) were melted at 1000 °C in an alumina crucible inside a sealed quartz tube, cooled down to 740 °C, and then cooled down to 610 °C at a rate of −1°C/h. Crystals were filtered in a centrifuge at this temperature. For sample B, Ca (Sigma-Aldrich, 99.99%) and Sb (Rare Metallic, 99.99%) were heated in a similar temperature profile but in a tungsten crucible inside a sealed quartz tube, cooled down to 740 °C, and then cooled down to 610 °C at a rate of −1°C/h. Crystals were filtered in a centrifuge at this temperature.

We found that typical size of the single-crystal domains is different between sample A (more than 1 mm size) and sample B (typically a few tens of μm). From the magnetic susceptibility measurements, we have confirmed the onset $T_c$ values of samples A and B to be ~1.8 K, as detailed in Fig. S1 in Sec. 1 of Ref. [41].

We performed ARPES measurements at BL28A in KEK-PF [42] and BL5U in UVSOR. We used circularly polarized photons of 30–300 eV at BL28A and linear horizontally polarized photons of 20–200 eV at BL5U. The energy and angular resolution was ~20 meV and 0.2°, respectively. The beam size on sample was 10×12 μm² [42] and 30×20 μm² for BL28A and BL5U, respectively. Samples were kept at 40 K during measurements. The Fermi level ($E_F$) of samples was referred to that of a gold film deposited on the sample holder. Laue backscattering measurements were performed prior to ARPES measurements to determine the sample geometry. A clear Laue pattern was observed for sample A as shown in Fig. 1(e), confirming the high single crystallinity of sample (note that it was difficult to obtain a clear Laue pattern for sample B due to the mixture of multiple crystal domains). However, we could safely determine the sample orientation for a small area of sample B on which the micro-ARPES measurements were carried out. We focused the micro photon beam on it. The cleaving plane (perpendicular to the c* axis) according to our previous focused x-ray diffraction and ARPES data. We observed no signature of aging or contamination of sample surface during ARPES measurements.

We have performed the first-principles band structure calculations by using Quantum ESPRESSO [43] with generalized gradient approximation [44] and have included the SOC. The Perdew-Burke-Ernzerhof (PBE) generalized gradient approximation was adopted for the exchange-correlation functional. Lattice constants were referenced from the experimental values [33] and a k mesh of 18×19×10 was used in the calculations.

### III. RESULTS AND DISCUSSION

To characterize the sample surface cleaved in UHV, we at first carried out photoemission measurements in a wide energy range with photons of $hν = 90$ eV. The energy distribution curve (EDC) for sample A is shown Fig. 1(f). We observe two peaks at the binding energy ($E_B$) of 31.5 and 32.7 eV, which are attributed to the spin-orbit-split Sb 4d5/2 and 4d3/2 orbital, respectively. Single-peaked structure of each orbital suggests that the chemical potential ($μ$) is almost the same between the Sb1 and Sb2 orbitals. A weak peak at $E_B = 25$ eV is assigned to the Ca 3p orbital, as better visualized in the expanded scale in the inset. No other peaks are observed in this $E_B$ range, supporting the clean surface.

Since CaSb₂ is a 3D material, it is essential to elucidate the bulk-band structure in the 3D BZ. For this sake, we have carried out $hν$-dependent ARPES measurements at the normal emission setup to observe the band dispersion along $k_z$ (parallel to the c* axis). We at first estimated the inner potential value to be $V_0 = 13.7$ eV from the periodicity of the band dispersion (for details, see Fig. 2). Then, we performed ARPES measurements at $hν = 118$ eV which corresponds to $k_z \sim 0$ plane, along the $\Gamma Y (k_z)$ cut at $T = 40$ K. The $\Gamma Y$ cut is expected to be useful to simultaneously observe the holelike band forming a 3D pocket at $\Gamma$ and an electronlike band forming a quasi-2D pocket at $\Gamma$, as suggested from the band calculation displayed in Figs. 1(c) and 1(d). As shown in Fig. 1(g), one can identify a highly dispersive holelike band at $E_B = 0–1$ eV which apparently crosses $E_F$ at $k_z = ± 0.18$ Å⁻¹ to form a hole pocket. According to the band calculation [36], this band originates from the Sb2 orbital [see Fig. 1(a)]. Inside this band, another holelike feature is seen at $E_B = 0.5–1$ eV. We found that the energy dispersion of these bands shows a qualitatively good agreement with the calculation along the $\Gamma Y$ cut shown by red curves in Fig. 1(h), whereas the calculated bands need to be shifted upward as a whole by 210 meV to obtain a best matching in the position of Fermi wave vectors ($k_F$’s) with the experiment. Surprisingly, we do not observe any signature of electron pockets at the $Y$ point originating from the Sb1 orbital [see Fig. 1(a)] [36], in sharp contrast to the calculation which does not take into account the chemical-potential ($μ$) shift [note that the $E_F$ position without the $μ$ shift ($E_F'$) is indicated by a dashed line in Fig. 1(h)]. This suggests that sample A is hole-doped probably due to the off-stoichiometry of the sample, leading to the lift-up of the 2D cylindrical electron bands into the unoccupied region. We have carefully surveyed the electron pocket at different in-plane ($k_x$ and $k_y$) and out-of-plane wave vectors ($k_z$) and found no evidence for it in the entire 3D BZ as detailed in Fig. 2. Therefore the transport property of this sample is expected to be governed by the 3D hole pocket, which is schematically illustrated by the calculated FS plot including the hole-doping effect in the bottom panel of Fig. 1(d).

The absence of the electron pocket is suggested from the FS mapping in the $k_x$ - $k_y$ plane [Fig. 2(b)] obtained by changing $hν$’s. Although the band calculation without the $μ$ shift shown in Fig. 2(c) (blue curves) predicted the existence of fairly vertical (i.e. quasi-2D) FSs elongated along the CY cut ($k_z \sim 0.75$ Å⁻¹), the corresponding ARPES intensity is totally absent as seen in Fig. 2(b). In the experiment, one can see a strong intensity surrounding the top $\Gamma$ point (at $k_z = 5.76$ Å⁻¹) which produces an ellipsoidal pocket elongated along the $\Gamma$Z direction (dashed red curve). This pocket is associated with the holelike band as seen from the ARPES intensity along the $k$ cut crossing the top $\Gamma$ point (cut 6) measured at $hν = 118$ eV in Fig. 2(d). On decreasing $hν$ (from cut 6 to cut 4), the separation between two $k_F$ points indicated by black arrows gradually becomes narrower due to the downward shift of the holelike band. Eventually the $k_F$ points disappear along
The gray rectangles indicate the calculation without the experimental hole pocket is much larger than that in the FS mapping obtained with different μ dependence of experimental band dispersions measured along the k cuts (1–6) shown in (a). ARPES data along cuts 1 and 6 nearly trace the same k region in reduced BZ but measured with different hv’s (118 and 88 eV). (e) In-plane FS mapping obtained with different hv’s of 118 eV (k_z ∼ 0; cut 7), 110 eV (k_z ∼ −0.5π; cut 8), and 103 eV (k_z ∼ −π; cut 9).

To clarify to what extent the observed band structure and fermiology is the intrinsic nature of CaSb_2 crystal, we carried out a similar ARPES measurement for sample B. Figure 3(a) shows the in-plane FS mapping at T = 40 K measured at hv = 90 eV which corresponds to the k_z ∼ 0 plane. One can see a dominant intensity around the Γ point elongated along the k_y direction. This feature originates from the Sb2-orbital-derived 3D hole pocket [36] as in the case of sample A, but its intensity distribution is different from that of sample A at k_z ∼ 0 shown in the bottom panel (slice 7) of Fig. 2(e), likely due to the difference in their doping levels as detailed below. The existence of a hole pocket is recognized from the clear E_F crossing of holelike band along a representative k cut (cut A) shown in Fig. 3(b). Besides the hole pocket, one can see in Fig. 3(a) a weak intensity around the Y point which is well separated from the hole pocket. Intriguingly, this feature is absent in sample A as seen in Figs. 2(b) and 2(e). The ARPES intensity along k cut crossing this feature (cut B) in Fig. 3(c) signifies the existence of a shallow electron pocket around the Y point. The metallic electron pocket is also confirmed by the energy distribution curve (EDC) at the Y point shown in Fig. 3(d) in which a peak near E_F associated with the electron-band bottom accompanied with the Fermi edge cutoff is clearly seen. We have confirmed the existence of electron pocket also at hv = 120 eV, in line with its quasi-2D nature predicted from the calculation shown in Figs. 1(c) and 1(d) (for details, see Sec. 2 of Ref. [41]). These results indicate that sample B is relatively more electron doped than sample A, and is situated in a semimetallic phase with both hole and electron pockets as predicted by the calculation for stoichiometric CaSb_2 (i.e., the calculation without μ shift). It is thus suggested that the absence of electron pocket in sample A is not due to the experimental artifacts such as the strong intensity suppression associated with the matrix-element effect but due to the difference in the doping levels. This conclusion is

![Diagram](https://via.placeholder.com/53x564.png)

FIG. 2. (a) Bulk BZ of CaSb_2 with the measured k cuts and k planes. Yellow rectangle represents the k plane (k_x-k_z plane) where the hv-dependent out-of-plane FS mapping shown in (b) was performed. Red curves indicate the k cuts (1–6) where the ARPES measurement shown in (d) was carried out. The gray rectangles indicate the k slices where in-plane FS mappings shown in (e) were performed. The slices 7, 8, and 9 correspond to k_z ∼ 0, −0.5π, and −π, respectively. (b) Out-of-plane FS mapping in the k_x-k_z plane obtained by sweeping hv. Dashed red curves are guide for the eyes to trace the Γ-centered pocket. (c) Corresponding calculated FSs with (red) and without (blue) taking into account the downward μ shift of 210 meV. (d) hv dependence of experimental band dispersions measured along the k cuts (1–6) shown in (a).
...the main player of superconductivity is the hole carriers in this pocket. It is thus inferred that the electron carriers in the Sb1-derived quasi-2D FS centered at the Y (C) point is not essential for the occurrence of superconductivity. This argument puts a constraint on the microscopic origin of superconductivity as well as its possible nontrivial nature in CaSb2, because the occupancy of the electron band hosting the line nodes was suggested to be important for promoting topological superconductivity associated with the dominant interorbital pairing interaction with odd-parity pairing [33]. Also, it is worth noting that CaSb2 maintains superconductivity regardless of the existence or absence of the electron pocket, which points to an intriguing possibility that one can tune the topological nature of superconductivity by simply controlling the carrier concentration in CaSb2, although its validation needs further experimental studies that directly connects the fermiology and the superconducting pairing symmetry. We leave such an experiment as a challenge in future.

**IV. CONCLUSION**

We reported results of a photon-energy-tunable micro-focused ARPES study on a topological line-nodal material CaSb2. We found two different types of FSs; one is a three-dimensional hole pocket at the Γ point and another is a cylindrical electron pocket at the Y (C) point which is predicted to host line nodes. When the doping level is changed to a more hole-doped region, the electron pocket at the Y (C) point disappears while the hole pocket at the Γ point survives, although the superconductivity keeps to emerge regardless of the doping level. This suggests a dominant role of the 3D hole pocket at the Γ point for the occurrence of superconductivity. The present result lays the foundation for understanding the nature of superconductivity and its relationship with topology in CaSb2.

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