Non-thermal hot electrons ultrafastly generating hot optical phonons in graphite

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Investigation of the non-equilibrium dynamics after an impulsive impact provides insights into couplings among various excitations. A two-temperature model (TTM) is often a starting point to understand the coupled dynamics of electrons and lattice vibrations: the optical pulse primarily raises the electronic temperature $T_e$ while leaving the lattice temperature $T_l$ low; subsequently the hot electrons heat up the lattice until $T_e = T_l$ is reached. This temporal hierarchy owes to the assumption that the electron-electron scattering rate is much larger than the electron-phonon scattering rate. We report herein that the TTM scheme is seriously invalidated in semimetal graphite. Time-resolved photoemission spectroscopy (TrPES) of graphite reveals that fingerprints of coupled optical phonons (COPs) occur from the initial moments where $T_e$ is still not definable. Our study shows that ultrafast-and-efficient phonon generations occur beyond the TTM scheme, presumably associated to the long duration of the non-thermal electrons in graphite.

One of the most interesting observation in the ultrafast dynamics of graphite is that $T_e$ in the sub-picosecond temporal region stays somewhat low even though irradiated with an intense femtosecond optical pulse. This led to the picture that the electronic energy is quasi-instantaneously transferred to the COPs through strong electron-phonon couplings based on the TTM scheme. However, electron-phonon coupling constant is reported to be moderately small in graphite, making the mechanism of the ultrafast COP generation elusive. Subsequent studies also suggest the nearly instantaneous COP generation coupled to the electron dynamics, which is considered to affect ballistic transports at high fields. Nevertheless, direct observation of the electron distribution in the transient is limited, and moreover, simultaneous detection of the electron distribution and the phonons in the transient has been beyond reach. TrPES [Fig. 1(a)] is one of the most powerful tools to investigate the dynamics of the electrons, since it can provide information of the transient electron distributions in a wide energy range across the Fermi level ($E_F$). Furthermore, Liu et al. recently reported that fingerprints of COPs occur in the photoemission spectra of graphite, as we shall explain below. Therefore, it became possible to monitor simultaneously the electron distribution and the fingerprints of COPs during the ultrafast dynamics of graphite by TrPES.

The fingerprints of COPs show up in the spectra recorded in a normal-emission geometry, that is, when we detect the photoelectrons around $\Gamma$ of the surface Brillouin zone of graphite [inset in Fig. 1(b)]. Since there are no bands around $\Gamma$ in the vicinity of $E_F$, the signal consists of photoelectrons around $K$ ($K'$) indirectly scattered into the vicinity of $\Gamma$ mainly by phonons. As we shall see later, signals of direct photoexcitations around the $K$ ($K'$) point indeed occur in the TrPES spectra recorded in the normal-emission geometry. A gap-like feature of $\sim 70$ meV occurs in the spectra near $E_F$ at low $T$ [Fig. 1(b)], since phonon absorption process is quenched and phonon emission (energyloss) process dominates. With increasing $T$, phonon absorption process is increased, and the spectral weight tails into higher energies resulting in an increase of the spectral intensity at $E_F$ [Fig. 1(b)]. This is similar to anti-Stokes lines in Raman spectra gaining stronger intensity at higher temperatures. Thus, we utilize the spectral intensity at $E_F$ in TrPES as a measure of the number of COPs in the transient. The optical phonons monitored herein are assigned to the 67-meV out-of-plane and the $\sim 150$-meV in-plane modes.
Results

TrPES spectra. Figure 2(a–f) show TrPES spectra of graphite, I(ω, t), recorded under a pump power p = 17 mW (a fluence of ~14 μJ/cm²) at room temperature. A movie file is provided in Supplementary Information. Overall, we observe that the electrons are pumped from the occupied side to the unoccupied side and subsequent recovery dynamics lasting over several tens of picoseconds: at t = 0 ps, the spectral intensity I(ω, t) is increased in the unoccupied side; at 0 ≤ t ≤ 0.20 ps, I(ω, t) at ω ~ 0.75 eV is decreased, whereas that at ω ~ 0 eV is increased; at t > 0.20 ps, I(ω, t) in the unoccupied side is decreased. Here, t = 0 ps has been determined utilizing the fast response of ~20 fs observed at ω = 0.8 eV, and the time resolution [full width at half the maximum (FWHM)] is estimated to be Δt = 0.43 ps, see Fig. 2(g).

We observe a plateau feature in the unoccupied side during the pump [Fig. 2(b) and the movie file in Supplementary Information], which is a hallmark of a non-thermal electron distribution. The plateau turns over into an exponential tailing at t > 0.20 ps [Fig. 2(f)], indicating that the electrons are distributed according to the Fermi-Dirac function. That is, electronic thermalization occurs at τ_e ≈ 0.2 ps so that T_ν becomes definable thereafter. In metallic materials, typical time scale for electronic thermalization is considered to be τ_e ~ 10 fs or less, and if Δt ≫ τ_e, one would not expect to observe a non-thermal distribution of the hot electrons. In fact, in the TrPES study of a metallic Bi_2Sr_2CaCu_2O_8, Perfetti et al. observed with Δt ~ 90 fs that the spectra mostly obey Fermi-Dirac statistics even at t ~ 0 ps. The smallness of 1/τ_e in graphite can be attributed to the semimetallic band structure: Since electronic states near E_F occur only around the K (K̄) points of the surface Brillouin zone, the available phase space for electron-electron scatterings becomes vanishingly small in approaching E_F particularly after the initial avalanche of the hot electrons towards E_F. This can act as a bottleneck for the non-thermal electrons to relax into the thermal (Fermi-Dirac) distribution.

The plateau observed just after the pump extends up to a cutoff at ω ~ 0.75 eV, as indicated by a bar in Fig. 2(b). This cutoff can be understood as a fingerprint of direct photoexcitations occurring around the K(K̄) point: since the π- and π*-bands are nearly symmetric about E_F as shown in the schematic in Fig. 2, direct excitations are dominated by the transitions from ω_hν/2 to ω_hν/2, and therefore, the cutoff energy can be identified to ω_hν/2 = 0.75 eV. Note that such a fingerprint of direct excitations would not be...
detected if $\tau_e \ll \Delta t$. Therefore, the $\sim 0.75$-eV cutoff strengthens the conclusion that the duration of the non-thermal electron distribution is detected in the present study.

With increasing $p$, the cutoff is blurred [inset in Fig. 2(d)], so that the unoccupied side of the spectra becomes featureless and the line shape becomes similar to the TrPES spectra of graphite reported previously.\textsuperscript{21} The intensity above the cutoff at $t = 0$ ps grows quadratically with $p$ while the intensity in the unoccupied side grows linearly with $p$ [Fig. 2(h)]. The latter indicates that the number of electrons excited by the pump is proportional to $p$ at least up to 78 mW, whereas the former indicates that the excited electrons are further scattered above the cutoff by other excitations such as the hot electrons themselves or hot phonons, since such scatterings occur roughly proportional to the square of the population of the excited particles. The intensity and the sharpness of the cutoff are almost independent of $T$ [inset of Fig. 2(f)], indicating that the blurring of the cutoff is not related to the thermally populated excitations in the initial state.

We now turn to the variation of the spectral weight around $E_F$, which serves as the fingerprints of COPs in the transient, as explained previously. One can see that the variation starts from the beginning of the transient in accord with the reports of the nearly instantaneous generation of the COPs in graphite.\textsuperscript{8, 11, 14} Our findings therefore show that the ultrafast COP generation takes place from the temporal region where the hot electrons are not thermalized [see, Fig. 2(a) and the movie file in Supplementary Information]. After $t \sim 0$ ps, the intensity at $E_F$ starts to decrease, see Fig. 2(e), indicating that the COP generation is mostly accomplished within $t \lesssim 0.2$ ps. The results are in strong contrast to the TTM scheme, where COPs are assumed to be cool at the beginning and then gradually heated up by the electrons that are thermalized.

Analysis of TrPES spectra. First, we investigate how the total electronic energy dissipates with time. As $I(\omega, t)$ reflects the occupied density of states (DOS), $\int_{0}^{\infty} \omega I(\omega, t) d\omega$ is a measure of the excess electronic energy, and we plot this as a function of $t$ in Fig. 3(a). Here, $\Delta I(\omega, t) = I(\omega, t) - I(\omega, -5 \text{ ps})$. The energy dissipation occurs with two time scales having distinct pump-power dependences: at $t = 1$ ps, the energy dissipation rate positively depends on $p$, whereas at $t > 1$ ps, it is independent of $p$. The fast dynamics at $t \lesssim 1$ ps is attributed to the net energy flow from the hot electrons to the hot COPs through electron-phonon scatterings, since this channel is increased quadratically with $p$ due to the increased populations of the hot electrons and hot COPs. Note that the scatterings among the hot electrons cannot account for the loss of the electronic
energy. At \( t \approx 1 \) ps, the electrons and COPs reach quasi-equilibrium (\( T_\text{eff} = T_{\text{COP}} \)) after a sufficient number of scatterings, and the electron-COP composite dissipates its energy to the heat bath such as acoustic phonons, and hence the \( p \)-independent energy dissipation at \( t \approx 1 \) ps.

Further analysis of the spectra also supports that the recovery dynamics at \( t \approx 1 \) ps is characterized by \( T_\text{eff} = T_{\text{COP}} \) and decay rates independent of \( p \). First, we derive the decay-rate spectrum for the slower component \( 1/\tau_2(\omega) \), which is obtained by fitting \( \tilde{f}(\omega, t) \) at each energy with a double exponential function (Supplementary Information). As shown in Fig. 3(b), \( 1/\tau_2(\omega) \) is independent of \( p \) and is quasi-linear to \( \omega \) at \( \omega \approx 0.3 \) eV. Second, we quantify the spectral shape by simulating the spectrum as \( \tilde{I}(\omega, t) = \alpha(t) \int G(\omega - \omega', \alpha^2(t) + \sigma_f^2) D(\omega') f(\omega', T_{\text{COP}}(t)) \, d\omega' \). Here, \( G(\omega, \alpha^2 + \sigma_f^2) \) is a Gaussian with FWHM of \( \sqrt{\alpha^2 + \sigma_f^2} \), \( f(\omega, T) \) is the Fermi-Dirac function, \( \alpha(t) \) is a scaling factor, and \( D(\omega) \) is a DOS. The Gaussian broadening accounts for the spectral weight accumulating from lower energies, so that \( \sigma_f \) becomes a measure of the number of the COPs in the transient. The spectra are nicely reproduced throughout the transient [Fig. 4(a) and Supplementary Information], and the fitting parameters are summarized in Fig. 4(b) and (c). One can see that \( \alpha^2 \) and \( \alpha \) for \( t > 1 \) ps scales with \( T_\text{COP} \) and does not explicitly depend on \( p \). This indicates that the number of the COPs at \( t > 1 \) ps is a function of \( T_{\text{COP}} \), i.e., \( T_{\text{COP}} = T_{\text{COP}} \) and that the spectral shape at \( t \approx 1 \) ps is determined only by \( T_{\text{COP}} \), which is the temperature of the electron-COP composite.

**Discussion**

We experimentally find that the electronic distribution is nonthermal at \( t \leq 0.2 \) ps, and also find evidence that the COP generation is mostly accomplished within this initial temporal region. This strongly indicates that the mechanism of the ultrafast COP generation is beyond the TTM scheme, see Fig. 5. Since the duration of the non-thermal electron distribution is longer than that in the TTM scheme, the high-energy electrons, which is considered to be favorable for generating high-energy COPs, indeed have more chance to generate COPs before they degrade into low-energy thermalized electrons. Therefore, not only the largeness of the electron-phonon couplings but also the long durations of the non-thermal electron distribution may be crucial for understanding the ultrafast-and-efficient generation of the COPs in graphite. It is interesting to note that a break-down of TTM was also suggested in \( \text{YBa}_2\text{Cu}_3\text{O}_7-\delta \), in which the DOS becomes vanishingly small around \( E_F \) in the \( d \)-wave superconducting state, similar to the case for graphite. Therefore, coupled dynamics directly involving the non-thermal electrons may be sought in materials that have vanishingly small DOS around \( E_F \) for example in neutral graphene, in nodal superconductors, and on the surface of topological insulators. Alternatively, we do not exclude the possibility that the hot COPs and hot electrons are co-generated by the pump, which may be viewed as a counterpart of the breakdown of adiabatic Born-Oppeenheimer approximation in as much as the electrons cannot follow the motion of the lattice, COPs are simultaneously generated when the electronic excitation takes place. Whichever the case may be, our study reveals that there is a unique mechanism of ultrafast COP generation where the concept of temperatures is broken.

**Methods**

The TrPES apparatus consists of an amplified Ti:sapphire laser system delivering \( h\nu = 1.5 \) eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into \( 1.5 \) eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into 1.5 eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into 1.5 eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into 1.5 eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into 1.5 eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into 1.5 eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a). A portion of the laser is converted into 1.5 eV pulses of 170-fs duration with 250-kHz repetition and a hemispherical analyzer, see Fig. 1(a).
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Author Contributions Statement
Y.I., T. Togashi, K.Y., M.T. and T. Taniuchi performed the experiments. Y.I. analyzed the data and wrote the paper with S.S. All the authors contributed to discussion.

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