Observation of charmless baryonic decays $B_{(s)}^0 \rightarrow p\bar{p}h^+h'^-$

The LHCb collaboration†

Abstract

Decays of $B^0$ and $B^0_s$ mesons to the charmless baryonic final states $p\bar{p}h^+h'^-$, where $h$ and $h'$ each denote a kaon or a pion, are searched for using the LHCb detector. The analysis is based on a sample of proton-proton collision data collected at center-of-mass energies of 7 and 8 TeV, corresponding to an integrated luminosity of 3 fb$^{-1}$. Four-body charmless baryonic $B_{(s)}^0$ decays are observed for the first time. The decays $B_{(s)}^0 \rightarrow p\bar{p}K^+K^-$, $B_{(s)}^0 \rightarrow p\bar{p}K^\pm\pi^\mp$, $B^0 \rightarrow p\bar{p}K^\pm\pi^\mp$ and $B^0 \rightarrow p\bar{p}\pi^+\pi^-$ are observed with a significance greater than 5 standard deviations; evidence at 4.1 standard deviations is found for the $B^0 \rightarrow p\bar{p}K^+K^-$ decay and an upper limit is set on the branching fraction for $B^0_s \rightarrow p\bar{p}\pi^+\pi^-$. Branching fractions in the kinematic region $m(p\bar{p}) < 2850$ MeV/c$^2$ are measured relative to the $B^0 \rightarrow J/\psi(\rightarrow p\bar{p})K^*(892)^0$ channel.

Published in Phys. Rev. D 96 (2017) 051103(R)

© CERN on behalf of the LHCb collaboration, licence [CC-BY-4.0](https://creativecommons.org/licenses/by/4.0/)

†Authors are listed at the end of this article.
In recent years, studies by the LHCb collaboration have greatly increased the knowledge of the decays of $B$ mesons to final states containing baryons. The first observation of a baryonic $B^+_c$ decay was reported in 2014 \cite{1}, and LHCb recently reported the first observation of a baryonic $B^+_s$ decay \cite{2}, the last of the four $B$ meson species for which a baryonic decay mode had yet to be observed.

Primary areas of interest in baryonic $B$ decays include the hierarchy of branching fractions to the various decay modes, the presence of resonances and the existence of a threshold enhancement in the baryon-antibaryon mass spectrum \cite{3,4}. First evidence of $CP$ violation in baryonic $B$ decays has been reported from an analysis of $B^+ \to p\bar{p}K^+$ decays \cite{5}. It is of great interest to search for further manifestations of $CP$ violation in baryonic $B$ decays, e.g. with so-called triple-product correlations (TPCs), see Ref. \cite{6} and references therein. For certain decays asymmetries of up to 20% are predicted \cite{7}. Four-body decays are particularly suited for this approach since the definition of the TPCs do not involve the spins of the final-state particles, unlike TPCs in three-body decays \cite{6,8}.

This paper presents a search for the decays of $B^0$ and $B^+_s$ mesons to the four-body charmless baryonic final states $p\bar{p}h^+h'^-$, where $h$ and $h'$ each denote a kaon or a pion. The inclusion of charge-conjugate processes is implied, unless otherwise indicated. For simplicity, the charges of the $h^+h'^-$ combinations will be omitted unless necessary. The branching fractions of these baryonic decays are measured relative to the $B^0 \to J/\psi(\to p\bar{p})K^+(892)^0$ channel. So far only the resonant decay $B^0 \to p\bar{p}K^+(892)^0$ has been seen by the BaBar \cite{9} and Belle \cite{10} collaborations, which measured its branching fraction to be $B(B^0 \to p\bar{p}K^+(892)^0) = (1.24^{+0.28}_{-0.25}) \times 10^{-6}$ \cite{11}. An upper limit $B(B^0 \to p\bar{p}\pi^+\pi^-) < 2.5 \times 10^{-4}$ at 90% confidence level has been set by the CLEO collaboration \cite{12}.

The data sample analyzed corresponds to an integrated luminosity of 1 fb$^{-1}$ of proton-proton collision data at a center-of-mass energy of 7 TeV and 2 fb$^{-1}$ at 8 TeV. The LHCb detector \cite{13,14} is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing $b$ or $c$ quarks. The detector elements that are particularly relevant to this analysis are as follows: a silicon-strip vertex detector surrounding the proton-proton interaction region that allows $c$ and $b$ hadrons to be identified from their characteristically long flight distance; a tracking system that provides a measurement of momentum, $p_T$, of charged particles; two ring-imaging Cherenkov detectors that are able to discriminate between different species of charged hadrons; and calorimeter and muon systems for the measurement of photons and neutral hadrons, and the detection of penetrating charged particles. Simulated data samples, produced with software described in Refs. \cite{15-20}, are used to evaluate the response of the detector and to investigate possible sources of background.

Real-time event selection is performed by a trigger \cite{21} that consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage, which performs a full event reconstruction. The hardware trigger stage requires events to have a muon with high transverse momentum, $p_T$, or a hadron, photon or electron with high transverse energy in the calorimeters. Signal candidates may come from events where the hardware trigger was caused either by signal particles or by other particles in the event. The software trigger requires a two-, three- or four-track secondary vertex with a significant displacement from any primary proton-proton interaction vertices (PVs). At least one charged particle must have $p_T > 1.6$ GeV/c and be inconsistent with originating from a PV. A multivariate algorithm \cite{22} is used for the identification of
secondary vertices consistent with the decay of a $b$ hadron.

The final selection of $B_{s}^{0}$ candidates, formed by combining four charged hadron candidates – a proton, an antiproton and an oppositely charged pair of light mesons – is carried out with a filtering stage followed by requirements on the response of a boosted decision tree (BDT) classifier\cite{23,24} and on particle identification (PID). The filtering stage includes requirements on the quality, $p$, $p_{T}$ and $\chi^{2}_{IP}$ of the tracks, loose PID requirements and an upper limit on the $p\bar{p}$ invariant mass; the $\chi^{2}_{IP}$ is defined as the difference between the vertex-fit $\chi^{2}$ of a PV reconstructed with and without the track in question. Each $B_{s}^{0}$ candidate must have a good-quality vertex that is displaced from the associated PV (that with which it forms the smallest $\chi^{2}_{IP}$), must satisfy $p$ and $p_{T}$ requirements, and must have a reconstructed invariant mass close to that of a $B_{s}^{0}$ meson under the signal mass hypothesis. A requirement is also imposed on the angle $\theta_{\text{dir}}$ between the candidate momentum vector and the line between the associated PV and the candidate decay vertex.

There are 15 input quantities to the BDT classifier: $p_{T}$, $\eta$, $\chi^{2}_{IP}$, $\theta_{\text{dir}}$ and the flight distance of the $B_{s}^{0}$ candidate; the quality of the $B_{s}^{0}$ vertex fit; the $p_{T}$ and $\chi^{2}_{IP}$ of the tracks; and the largest distance of closest approach between any pair of tracks. The BDT is trained using simulated $B_{s}^{0} \rightarrow p\bar{p}hh'$ signal candidates, generated with uniform distributions over phase space, and events in a high sideband of the $p\bar{p}K\pi$ invariant mass in data ($m(p\bar{p}K\pi)$ in the range 5450–5550 MeV/$c^{2}$) to represent the background. Tight PID requirements are applied to all final-state particles to reduce the combinatorial background, suppress the cross-feed backgrounds between the different $p\bar{p}hh'$ final states — background from other signal decays where one particle is misidentified — and ensure that the datasets for the three $p\bar{p}hh'$ final states are mutually exclusive. For each final state individually, the requirements on the PID and BDT response are optimized for the signal significance using simulation samples for the signal. After all selection requirements are applied, approximately 3% of events with at least one candidate also contain a second candidate; a candidate is then selected at random. The efficiency of the full reconstruction and selection, including the acceptance and the trigger selection, is approximately 0.1%.

To reject contributions from intermediate charm states, candidates with $hh'$ invariant mass consistent with a $D^{0}$ meson or $phh'$ invariant mass consistent with a $\Lambda^{+}_{c}$ baryon are removed. The contribution from the charmonium region is removed by requiring the invariant mass of the $p\bar{p}$ pair to be less than 2850 MeV/$c^{2}$, similar to the procedure in Refs.\cite{5,25}. This last requirement is not applied to the normalization mode $B^{0} \rightarrow J/\psi K^{*}(892)^{0}$, where the vector mesons are reconstructed in the $J/\psi \rightarrow p\bar{p}$ and $K^{*}(892)^{0} \rightarrow K^{+}\pi^{-}$ decay modes. All the other steps of the selection are in common for the signal and the normalization modes.

The yields of the signal decays are obtained from a simultaneous unbinned extended maximum likelihood fit to the $B_{s}^{0}$ candidate invariant mass distributions in the three $p\bar{p}hh'$ final states in the range 5165–5525 MeV/$c^{2}$. This approach accounts for potential cross-feed from one channel to another due to particle misidentification. Each signal component is modeled with a double-sided Crystal Ball (DSCB) function\cite{26}. For each signal the tail parameters of the DSCB functions are determined from simulation. The peak position of the $B^{0}$ signals is common to the three final states, while the difference between the peak positions of the $B^{0}$ and $B_{s}^{0}$ signals is constrained to its known value\cite{11}. The width of the $B^{0}$ signal is a free parameter in the $p\bar{p}K\pi$ final state and it is related to the width in the other two final states by scale factors determined from simulation. The
same applies to the width of the $B_s^0$ signals, which is a free parameter only in the $p\bar{p}KK$ final state.

For each final state the dominant $B_s^0 \to p\bar{p}hh'$ cross-feed background is included: the $B^0 \to p\bar{p}K\pi$ mode in the $p\bar{p}KK$ and $p\bar{p}\pi\pi$ invariant mass distributions, and the $B^0 \to p\bar{p}\pi\pi$ mode in the $p\bar{p}K\pi$ spectrum. Each cross-feed background is modeled with a DSCB function with all the shape parameters fixed according to simulation; the yield is fixed relative to the yield in the correctly reconstructed final state taking into account the (mis)identification probabilities calibrated using data, as described below. In addition, a combinatorial background component modeled by an exponential function, with both parameters free to vary, is present for each final state.

The yield of the normalization decay is determined from a separate simultaneous fit to the $p\bar{p}K\pi$, $p\bar{p}$ and $K\pi$ invariant mass distributions in the ranges 5180–5380 MeV/$c^2$, 3047–3147 MeV/$c^2$ and 642–1092 MeV/$c^2$, respectively. The $B^0 \to J/\psi K^*(892)^0$ component is parameterized in the $K\pi$ invariant mass distribution by a relativistic spin-1 Breit-Wigner function and in the $p\bar{p}K\pi$ and $p\bar{p}$ invariant mass distributions by DSCB functions with the tail parameters fixed from simulation. The $K\pi$ S-wave component is modeled in the $K\pi$ invariant mass distribution by the LASS parametrization \cite{27,28} that describes nonresonant and $K^*_{s}(1430)^0$ S-wave contributions; this component is modeled in the $p\bar{p}K\pi$ and $p\bar{p}$ invariant mass distributions with the same shape as the $B^0 \to J/\psi K^*(892)^0$ component. A combinatorial background component modeled by a freely varying exponential function is also present in each spectrum.

The $p\bar{p}hh'$ invariant mass distributions with the results of the fit overlaid are shown in Fig. 1 while the signal yields and the significances are collected in Table 1. The significance of each of the signal modes is determined from the change in likelihood in Fig. 1 while the signal yields and the significances are collected in Table 1. The $K\pi\pi$ invariant mass distributions in the ranges 5180–5380 MeV/$c^2$, 3047–3147 MeV/$c^2$ and 642–1092 MeV/$c^2$, respectively, while the other signal modes have significances greater than 25 $\sigma$.

The branching fractions of the $B_s^0 \to p\bar{p}hh'$ decays are determined relative to the visible branching fraction of the $B^0 \to J/\psi K^*(892)^0$ decay using

$$\frac{B(B^0 \to J/\psi K^*(892)^0)}{B_{\text{vis}}(B^0 \to J/\psi K^*(892)^0)} = \frac{N^\text{corr}(B^0_s \to p\bar{p}hh')}{N^\text{corr}(B^0 \to J/\psi K^*(892)^0)} \left( \frac{f_d}{f_s} \right),$$

where $f_s/f_d = 0.259 \pm 0.015$ (included only for the $B_s^0$) is the ratio of $b$ hadronization probabilities, $f_q$, to the hadron $B_q$ \cite{30}, and $N^\text{corr}$ denote efficiency-corrected fitted signal yields. The yields are obtained from the mass fits, while simulation is used to evaluate the contribution to the efficiency from each stage of the selection except for the effect of the PID criteria. The latter is determined from calibration data samples of kinematically identified pions, kaons and protons originating from the decays $D^{\ast+} \to D^0 (\to K^-\pi^+)\pi^+$, $\Lambda \to p\pi^−$ and $\Lambda_c^+ \to pK^-\pi^+$ and weighted according to the kinematics of the signal particles \cite{31,32}. For each final state the efficiencies are determined as a function of the position in phase space, and efficiency corrections for each candidate are applied using the method of Ref. \cite{33} to take the variation over the phase space into account. Explicitly, $N^\text{corr} = \sum_i W_i / \varepsilon_i$, where the sum runs over the candidates in the fit, $W_i$ is the $s$Weight for candidate $i$ determined with the $sPlot$ method \cite{34} and $\varepsilon_i$ is the efficiency for the candidate $i$ which depends only on its position in the five-dimensional phase space. The visible branching fraction of the normalization mode, defined as $B(B^0 \to J/\psi K^*(892)^0) \times$
Figure 1: Invariant mass distributions for $B^0_s$ candidates in the (top left) $p\bar{p}K'K$, (top right) $p\bar{p}KK$ final state and (bottom right) invariant mass distribution of $B^0 \rightarrow J/\psi K^*(892)^0$ in the $p\bar{p}K\pi$ final state. The results of the fits are shown with blue solid lines. In the first three figures signals for $B^0$ and $B^0_s$ decays are shown, respectively, with green dotted and red dot-dashed lines, combinatorial backgrounds are shown with black dashed lines and cross-feed backgrounds are shown with violet dot-dashed lines. In the bottom right figure the normalization signal is shown with a green dotted line, the $K\pi$ S-wave component is displayed with a red dot-dashed line and the combinatorial background with a black dashed line.

$B(J/\psi \rightarrow p\bar{p}) \times B(K^*(892)^0 \rightarrow K^+\pi^-)$, is $B_{\text{vis}}(B^0 \rightarrow J/\psi K^*(892)^0) = (1.68 \pm 0.12) \times 10^{-6}$, where the $B^0 \rightarrow J/\psi K^*(892)^0$ branching fraction is taken from Ref. 35 and the others from Ref. 11.

The branching fraction of each signal mode is reported in Table 1. The significance for the $B^0_s \rightarrow p\bar{p}\pi\pi$ mode is less than $3\sigma$; an upper limit on its branching fraction is found to be

$$B(B^0_s \rightarrow p\bar{p}\pi\pi) < 6.6 \times 10^{-7}$$

at 90% confidence level, by integrating the likelihood after multiplying by a prior probability distribution that is uniform in the region of positive branching fraction. The values of the ratios of branching fractions between different $B^0_s \rightarrow p\bar{p}hh'$ decay modes are reported in Table 2.

The signal distributions in $m(hh')$ and $m(p\bar{p})$ are obtained by subtracting the background using the sPlot technique 34, with the $B^0_s$ candidate invariant mass as the discriminating variable. Per-candidate weights are applied to correct for the variation of the selection efficiency over the phase space. Figure 2 shows the $hh'$ invariant mass distributions of the $B^0 \rightarrow p\bar{p}K\pi$, $B^0_s \rightarrow p\bar{p}KK$ and $B^0 \rightarrow p\bar{p}\pi\pi$ decay modes. A peak from a vector meson is identifiable in each mass spectrum, corresponding to a $K^*(892)^0$, a
The uncertainties on the yields are statistical only. The first uncertainty on each branching fraction is statistical, the second systematic, the third comes from the uncertainty on the branching fraction of the normalization mode and the fourth, where present, is due to the uncertainty on $f_d/f_s$.

| Decay channel | Yield $N$ | Significance $[\sigma]$ | Branching fraction / $10^{-6}$ |
|---------------|-----------|--------------------------|-------------------------------|
| $B^0 \to p\bar{p}KK$ | $68 \pm 17$ | 4.1 | $0.113 \pm 0.028 \pm 0.011 \pm 0.008$ |
| $B^0 \to p\bar{p}K\pi$ | $4155 \pm 83$ | $> 25$ | $5.9 \pm 0.3 \pm 0.3 \pm 0.4$ |
| $B^0 \to p\bar{p}\pi\pi$ | $902 \pm 35$ | $> 25$ | $2.7 \pm 0.1 \pm 0.1 \pm 0.2$ |
| $B^0_s \to p\bar{p}KK$ | $635 \pm 32$ | $> 25$ | $4.2 \pm 0.3 \pm 0.2 \pm 0.3 \pm 0.2$ |
| $B^0_s \to p\bar{p}K\pi$ | $246 \pm 39$ | 6.5 | $1.30 \pm 0.21 \pm 0.11 \pm 0.09 \pm 0.08$ |
| $B^0_s \to p\bar{p}\pi\pi$ | $39 \pm 16$ | 2.6 | $0.41 \pm 0.17 \pm 0.04 \pm 0.03 \pm 0.02$ |
| $B^0 \to J/\psi K^*(892)^0$ | $1216 \pm 45$ | – | – |

The sources of systematic uncertainty on the absolute branching fractions and on the ratios of branching fractions arise from the fit model, the knowledge of the efficiencies and, where appropriate, from the uncertainties on the branching fraction of the normalization mode and on the ratio of $b$-quark hadronization probabilities. Pseudoexperiments are used to estimate the effect of using alternative shapes for the fit components, or of including additional components in the fit. In particular, the effect of adding other cross-feed backgrounds, partially reconstructed backgrounds and components coming from $A_0^0$ decays have been investigated. These are the dominant sources of systematic uncertainty for the $B^0 \to p\bar{p}KK$ and $B^0 \to p\bar{p}\pi\pi$ modes. The effect of fixing the yields of the cross-feed backgrounds based on the (mis)identification probabilities is also assessed by varying these probabilities within their uncertainties. Intrinsic biases in the fitted yields are investigated with pseudoexperiments and are found to be negligible. Uncertainties on the efficiencies, due to the limited size of the simulation samples, the uncertainty on their evaluated distributions across the phase space of the decays and from possible residual differences between data and simulation. The unknown decay kinematics are the principal source of systematic uncertainty for the $B^0_s \to p\bar{p}K\pi$ mode, while for the $B^0 \to p\bar{p}KK$, $B^0 \to p\bar{p}K\pi$ and $B^0 \to p\bar{p}\pi\pi$ modes the dominant source of systematic uncertainty comes from the

Table 2: Ratios of branching fractions among different $B^{0}_{(s)} \to p\bar{p}hh'$ modes. The first uncertainty is statistical, the second systematic and the third, where present, comes from the uncertainty on $f_d/f_s$.

| Equation | Ratio | Uncertainty |
|----------|-------|-------------|
| $B(B^0 \to p\bar{p}KK)/B(B^0 \to p\bar{p}K\pi)$ | $0.019 \pm 0.005 \pm 0.002$ |
| $B(B^0_s \to p\bar{p}K\pi)/B(B^0 \to p\bar{p}K\pi)$ | $0.46 \pm 0.02 \pm 0.02$ |
| $B(B^0 \to p\bar{p}K\pi)/B(B^0 \to p\bar{p}K\pi)$ | $0.6 \pm 0.04 \pm 0.02 \pm 0.01$ |
| $B(B^0_s \to p\bar{p}K\pi)/B(B^0 \to p\bar{p}K\pi)$ | $0.31 \pm 0.05 \pm 0.02$ |

$\phi(1020)$ and a $\rho(770)^0$ meson, respectively. The $p\bar{p}$ invariant mass distributions are also shown for the same decay modes. An enhancement near threshold, typical in baryonic $B$ decays,$^3,4,$ is clearly visible in each case. Detailed amplitude analyses of the $B^{0}_{(s)} \to p\bar{p}hh'$ decays will be of interest with larger samples.
uncertainty on the efficiency of the hardware stage of the trigger. As the efficiencies depend on the signal decay-time distribution, the effect coming from the different lifetimes of the $B_s^0$ mass eigenstates has been evaluated. The systematic uncertainties due to the vetoes of charm hadrons are also included.

In summary, a search for the four-body charmless baryonic decays $B^0 \rightarrow p\bar{p}hh'$ has been carried out by the LHCb collaboration with a sample of proton-proton collision data corresponding to an integrated luminosity of 3 fb$^{-1}$. First observations are obtained for the decays $B^0 \rightarrow p\bar{p}\pi\pi$, nonresonant $B^0 \rightarrow p\bar{p}K\pi$, $B^0_s \rightarrow p\bar{p}KK$ and $B^0_s \rightarrow p\bar{p}K\pi$, while first evidence is reported for the $B^0 \rightarrow p\bar{p}KK$ mode and an upper limit is set on the
$B_s^0 \to p\bar{p}\pi\pi$ branching fraction. In particular, four-body baryonic $B_s^0$ decays are observed for the first time and a threshold enhancement in the baryon-antibaryon mass spectra is confirmed for baryonic $B_s^0$ decays [2].

The LHCb collaboration has recently published studies of $CP$ violation with four-body $\Lambda_b^0 \to p h^- h^+ h^-$ decays studying triple-product correlations, and presented first evidence for $CP$ violation in baryons [36]. The decays of $B^0$ and $B_s^0$ mesons to $p\bar{p}hh'$ final states reported in this paper may be used in the future for similar studies of $CP$ violation in baryonic $B$ decays.

Acknowledgements

We express our gratitude to our colleagues in the CERN accelerator departments for the excellent performance of the LHC. We thank the technical and administrative staff at the LHCb institutes. We acknowledge support from CERN and from the national agencies: CAPES, CNPq, FAPERJ and FINEP (Brazil); MOST and NSFC (China); CNRS/IN2P3 (France); BMBF, DFG and MPG (Germany); INFN (Italy); NWO (The Netherlands); MNiSW and NCN (Poland); MEN/IFA (Romania); MinES and FASO (Russia); MinECo (Spain); SNSF and SER (Switzerland); NASU (Ukraine); STFC (United Kingdom); NSF (USA). We acknowledge the computing resources that are provided by CERN, IN2P3 (France), KIT and DESY (Germany), INFN (Italy), SURF (The Netherlands), PIC (Spain), GridPP (United Kingdom), RRCKI and Yandex LLC (Russia), CSCS (Switzerland), IFIN-HH (Romania), CBPF (Brazil), PL-GRID (Poland) and OSC (USA). We are indebted to the communities behind the multiple open source software packages on which we depend. Individual groups or members have received support from AvH Foundation (Germany), EPLANET, Marie Sklodowska-Curie Actions and ERC (European Union), Conseil Général de Haute-Savoie, Labex ENIGMASS and OCEVU, Région Auvergne (France), RFBR and Yandex LLC (Russia), GVA, XuntaGal and GENCAT (Spain), Herchel Smith Fund, The Royal Society, Royal Commission for the Exhibition of 1851 and the Leverhulme Trust (United Kingdom).
References

[1] LHCb collaboration, R. Aaij et al., First observation of a baryonic $B^+_c$ decay, Phys. Rev. Lett. 113 (2014) 152003, arXiv:1408.0971.

[2] LHCb collaboration, R. Aaij et al., First observation of a baryonic $B^0_s$ decay, arXiv:1704.07908, submitted to Phys. Rev. Lett.

[3] W.-S. Hou and A. Soni, Pathways to rare baryonic $B$ decays, Phys. Rev. Lett. 86 (2001) 4247, arXiv:hep-ph/0008079.

[4] BaBar and Belle collaborations, A. J. Bevan et al., The physics of the $B$ factories, Eur. Phys. J. C74 (2014) 3026, arXiv:1406.6311.

[5] LHCb collaboration, R. Aaij et al., Evidence for $CP$ violation in $B^+ \rightarrow p\bar{p}K^+$ decays, Phys. Rev. Lett. 113 (2014) 141801, arXiv:1407.5907.

[6] C. Q. Geng and Y. K. Hsiao, Direct $CP$ and $T$ violation in baryonic $B$ decays, Int. J. Mod. Phys. A23 (2008) 3290, arXiv:0801.0022.

[7] C. Q. Geng, Y. K. Hsiao, and J. N. Ng, Direct $CP$ violation in $B^\pm \rightarrow p\bar{p}K^{(*)\pm}$, Phys. Rev. Lett. 98 (2007) 011801, arXiv:0608328.

[8] M. Gronau and J. L. Rosner, Triple product asymmetries in $K$, $D_{(s)}$ and $B_{(s)}$ decays, Phys. Rev. D84 (2011) 096013, arXiv:1107.1232.

[9] BaBar collaboration, B. Aubert et al., Evidence for the $B^0 \rightarrow p\bar{p}K^{*0}$ and $B^+ \rightarrow \eta_cK^{*+}$ decays and study of the decay dynamics of $B$ meson decays into $p\bar{p}h$ final states, Phys. Rev. D76 (2007) 092004, arXiv:0707.1648.

[10] Belle collaboration, J. H. Chen et al., Observation of $B^0 \rightarrow p\bar{p}K^{*0}$ with a large $K^{*0}$ polarization, Phys. Rev. Lett. 100 (2008) 251801, arXiv:0802.0336.

[11] Particle Data Group, C. Patrignani et al., Review of particle physics, Chin. Phys. C40 (2016) 100001.

[12] CLEO collaboration, C. Bebek et al., Search for the charmless decays $B \rightarrow p\bar{p}\pi$ and $p\bar{p}\pi\pi$, Phys. Rev. Lett. 62 (1989) 8.

[13] LHCb collaboration, A. A. Alves Jr. et al., The LHCb detector at the LHC, JINST 3 (2008) S08005.

[14] LHCb collaboration, R. Aaij et al., LHCb detector performance, Int. J. Mod. Phys. A30 (2015) 1530022, arXiv:1412.6352.

[15] T. Sjöstrand, S. Mrenna, and P. Skands, PYTHIA 6.4 physics and manual, JHEP 05 (2006) 026, arXiv:hep-ph/0603175. T. Sjöstrand, S. Mrenna, and P. Skands, A brief introduction to PYTHIA 8.1, Comput. Phys. Commun. 178 (2008) 852, arXiv:0710.3820.

[16] I. Belyaev et al., Handling of the generation of primary events in Gauss, the LHCb simulation framework, J. Phys. Conf. Ser. 331 (2011) 032047.
D. J. Lange, The EvtGen particle decay simulation package, Nucl. Instrum. Meth. A462 (2001) 152.

P. Golonka and Z. Was, PHOTOS Monte Carlo: A precision tool for QED corrections in Z and W decays, Eur. Phys. J. C45 (2006) 97, arXiv:hep-ph/0506026.

Geant4 collaboration, J. Allison et al., Geant4 developments and applications, IEEE Trans. Nucl. Sci. 53 (2006) 270; Geant4 collaboration, S. Agostinelli et al., Geant4: A simulation toolkit, Nucl. Instrum. Meth. A506 (2003) 250.

M. Clemencic et al., The LHCb simulation application, Gauss: Design, evolution and experience, J. Phys. Conf. Ser. 331 (2011) 032023.

R. Aaij et al., The LHCb trigger and its performance in 2011, JINST 8 (2013) P04022, arXiv:1211.3055.

V. V. Gligorov and M. Williams, Efficient, reliable and fast high-level triggering using a bonsai boosted decision tree, JINST 8 (2013) P02013, arXiv:1210.6861.

L. Breiman, J. H. Friedman, R. A. Olshen, and C. J. Stone, Classification and regression trees, Wadsworth international group, Belmont, California, USA, 1984.

Y. Freund and R. E. Schapire, A decision-theoretic generalization of on-line learning and an application to boosting, J. Comput. Syst. Sci. 55 (1997) 119.

LHCb collaboration, R. Aaij et al., Search for $B^+_c$ decays to the $p\bar{p}\pi^+$ final state, Phys. Lett. B759 (2016) 313, arXiv:1603.07037.

T. Skwarnicki, A study of the radiative cascade transitions between the Upsilon-prime and Upsilon resonances, PhD thesis, Institute of Nuclear Physics, Krakow, 1986, DESY-F31-86-02.

D. Aston et al., A study of $K^-\pi^+$ scattering in the reaction $K^-p \rightarrow K^-\pi^+n$ at 11 GeV/c, Nucl. Phys. B296 (1988) 493.

BaBar collaboration, B. Aubert et al., Time-dependent and time-integrated angular analysis of $B \rightarrow \phi K_s^0\pi^0$ and $\phi K^\pm\pi^\mp$, Phys. Rev. D78 (2008) 092008, arXiv:0808.3586.

S. S. Wilks, The large-sample distribution of the likelihood ratio for testing composite hypotheses, Ann. Math. Stat. 9 (1938) 60.

LHCb collaboration, R. Aaij et al., Measurement of the fragmentation fraction ratio $f_s/f_d$ and its dependence on B meson kinematics, JHEP 04 (2013) 001, arXiv:1301.5286; $f_s/f_d$ value updated in LHCb-CONF-2013-011.

M. Adinolfi et al., Performance of the LHCb RICH detector at the LHC, Eur. Phys. J. C73 (2013) 2431, arXiv:1211.6759.

L. Anderlini et al., The PIDCalib package, LHCb-PUB-2016-021.

LHCb collaboration, R. Aaij et al., Observation of $B^0 \rightarrow \bar{D}^0 K^+K^-$ and evidence for $B^0_s \rightarrow \bar{D}^0 K^+K^-$, Phys. Rev. Lett. 109 (2012) 131801, arXiv:1207.5991.
[34] M. Pivk and F. R. Le Diberder, *sPlot: A statistical tool to unfold data distributions*, Nucl. Instrum. Meth. A555 (2005) 356, arXiv:physics/0402083.

[35] Belle collaboration, K. Chilikin et al., *Observation of a new charged charmoniumlike state in $B^0 \rightarrow J/\psi K^-\pi^+$ decays*, Phys. Rev. D90 (2014) 112009, arXiv:1408.6457.

[36] LHCb collaboration, R. Aaij et al., *Measurement of matter-antimatter differences in beauty baryon decays*, Nature Physics 13 (2017) 391, arXiv:1609.05216.
LHCb collaboration

R. Aaij, B. Adeeva, M. Adinolfi, Z. Ajaltouni, S. Akar, J. Albrecht, F. Alessio, M. Alexander, S. Ali, G. Alkhazov, P. Alvarez Cartelle, A.A. Alves Jr, S. Amato, S. Amerio, Y. Ambhis, L. An, L. Anderlini, G. Andreassi, M. Andreotti, J.E. Andrews, R.B. Appleby, F. Archilli, P. d’Argent, J. Arnau Romeu, A. Artamonov, M. Artuso, E. Aslanides, G. Auriemma, M. Baalouch, I. Babuschkin, S. Bachmann, J.J. Back, A. Badalov, C. Baesso, S. Bakes, V. Balagura, W. Baldini, A. Baranov, R.J. Barlow, C. Barschel, S. Barsuk, W. Barter, F. Baryshnikov, M. Basiczek, B. Vatozkaya, B. Batsukh, V. Battista, A. Bay, L. Beaucourt, J. Beddow, F. Bedeschi, I. Bediaga, A. Beiter, L.J. Bel, V. Belli, N. Belloli, K. Belous, I. Belyaev, E. Ben-Haim, G. Bencivenni, S. Benson, B. Bernacki, A. Berezhnoy, R. Bernet, A. Bertoncini, C. Betancourt, F. Bettì, M.-O. Bétiller, M. van Beuzekom, I. Bezhshyo, S. Bifani, P. Billior, A. Birnkrant, A. Bitadze, A. Bizietz, T. Blanc, J. Blouw, S. Blusk, V. Bocci, T. Boettcher, A. Bondar, N. Bondar, W. Bonivento, I. Bordyuzhin, A. Borgheresi, B. Borgi, M. Borisayyk, M. Borsato, F. Bosso, M. Boudin, T.J.V. Bowcock, E. Bowen, C. Bozzi, S. Braun, T. Britton, J. Brodzicka, E. Buchanan, C. Burt, A. Bursche, J. Buytaert, S. Cadeddu, R. Calabrese, M. Calvi, M. Calvo, G. Camci, P. Campagna, D.H. Campos, L. Capriotti, A. Carboni, G. Carboni, R. Cardinale, A. Cardini, P. Carniti, L. Carson, K. Carvalho Akiba, G. Casser, L. Cassina, L. Castillo Garcia, M. Cattaneo, G. Cavallero, R. Cenci, D. Chamont, M. Charles, Ph. Charpentier, G. Chatzikostantinidis, M. Chedeville, S. Chen, S.F. Cheung, V. Chobanova, M. Chrzaszcz, A. Chubynkin, X. Cid Vidal, G. Ciezarek, P.E.L. Clarke, M. Clemenčič, H.V. Cliff, J. Closier, V. Coco, J. Cogan, E. Cogneras, V. Cogoni, L. Cojocariu, P. Collins, A. Comerma-Montells, A. Contu, A. Cook, G. Coombs, S. Coquereau, G. Corti, M. Corvo, C.M. Costa, Sobia, B. Couturier, G.A. Cowan, D.C. Craik, A. Crocombe, M. Cruz Torres, S. Cunliffe, R. Currie, C.D. D’Ambrosio, F. Da Cunha Marinho, E. Dall’Occo, J. Dalsero, A. Davis, K. De Bruyn, S. De Capua, M. De Cian, J.M. De Miranda, L. De Paula, M. De Serio, P. De Simone, A. Dean, K. Decamp, M. Deckenhoff, L. Del Buono, H.-P. Dembinski, M. Demmer, A. Dendek, D. Derkach, O. Deschamps, F. Dettori, B. Dey, A. Di Canto, P. Di Nezza, H. Diikstra, F. Dordoi, M. Dorigo, A. Dosil Suárez, A. Dovbnya, K. Dreimanis, L. Dufour, G. Dujany, K. Dungs, P. Durante, R. Dzhelyadin, M. Dziewiecki, A. Dziurda, A. Dzyuba, S. Déléage, S. Easo, M. Ebert, U. Egede, V. Egorychev, S. Eidelman, S. Eisenhardt, U. Eitschberger, R. Ekelhof, L. Eklund, S. Elly, S. Esen, H.M. Evans, T. Evans, A. Fabalba, N. Farley, S. Farry, R. Fay, D. Fazio, C.M. Fargion, G. Fernandez, A. Fernandez Prieto, F. Ferrari, F. Ferreira Rodrigues, M. Ferro-Luzzi, S. Filipov, R.A. Fini, M. Fiore, M. Fiorini, M. Firlej, C. Fitzpatrick, T. Fiantowski, F. Fleuret, K. Fohl, M. Fontana, F. Fontanelli, D.C. Forshaw, R. Forty, V. Franco Lima, M. Frank, C. Frei, J. Fu, W. Funk, E. Furfaro, C. Färber, A. Gallas Torreira, D. Galli, S. Gallorini, S. Gambetta, M. Gandelman, P. Gandini, Y. Gao, L.M. Garcia Martin, J. Garcia Pardina, J. Garra Tico, L. Garrido, P.J. Garsed, D. Gascon, C. Gaspar, L. Gavardi, G. Gazzoni, J. Gerick, E. Gersabeck, M. Gersabeck, T. Gershon, Ph. Ghez, S. Giani, V. Gibson, O.G. Girard, L. Giubega, K. Gizdov, V.V. Gilgorov, D. Golubkov, A. Golutvin, A. Gomes, I.V. Gorelov, C. Gotti, E. Govorkova, R. Graciani, D. Grazzini, L.A. Granado Cardoso, E. Grauges, E. Graverini, G. Graziani, A. Greco, R. Greim, P. Griffith, L. Grillo, B.R. Gruberg Cazon, O. Grünberg, E. Gushchin, Y. Guz,
Laboratori Nazionali dell’INFN di Frascati, Frascati, Italy
Sezione INFN di Genova, Genova, Italy
Università & INFN, Milano-Bicocca, Milano, Italy
Sezione di Milano, Milano, Italy
Sezione INFN di Padova, Padova, Italy
Sezione INFN di Pisa, Pisa, Italy
Sezione INFN di Roma Tor Vergata, Roma, Italy
Sezione INFN di Roma La Sapienza, Roma, Italy
Henryk Niewodniczanski Institute of Nuclear Physics Polish Academy of Sciences, Kraków, Poland
AGH - University of Science and Technology, Faculty of Physics and Applied Computer Science, Kraków, Poland
National Center for Nuclear Research (NCBJ), Warsaw, Poland
Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest-Magurele, Romania
Petersbury Nuclear Physics Institute (PNPI), Gatchina, Russia
Institute of Theoretical and Experimental Physics (ITEP), Moscow, Russia
Institute of Nuclear Physics, Moscow State University (SINP MSU), Moscow, Russia
Institute for Nuclear Research of the Russian Academy of Sciences (INR RAN), Moscow, Russia
Yandex School of Data Analysis, Moscow, Russia
Budker Institute of Nuclear Physics (SB RAS), Novosibirsk, Russia
Institute for High Energy Physics (IHEP), Protvino, Russia
ICCUB, Universitat de Barcelona, Barcelona, Spain
Universidad de Santiago de Compostela, Santiago de Compostela, Spain
European Organization for Nuclear Research (CERN), Geneva, Switzerland
Institute of Physics, Ecole Polytechnique Fédérale de Lausanne (EPFL), Lausanne, Switzerland
Physik-Institut, Universität Zürich, Zürich, Switzerland
Nikhef National Institute for Subatomic Physics, Amsterdam, The Netherlands
Nikhef National Institute for Subatomic Physics and VU University Amsterdam, Amsterdam, The Netherlands
NSC Kharkiv Institute of Physics and Technology (NSC KIPT), Kharkiv, Ukraine
Institute for Nuclear Research of the National Academy of Sciences (KINR), Kyiv, Ukraine
University of Birmingham, Birmingham, United Kingdom
H.H. Wills Physics Laboratory, University of Bristol, Bristol, United Kingdom
Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
Department of Physics, University of Warwick, Coventry, United Kingdom
STFC Rutherford Appleton Laboratory, Didcot, United Kingdom
School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
Imperial College London, London, United Kingdom
School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
Department of Physics, University of Oxford, Oxford, United Kingdom
Massachusetts Institute of Technology, Cambridge, MA, United States
University of Cincinnati, Cincinnati, OH, United States
University of Maryland, College Park, MD, United States
Syracuse University, Syracuse, NY, United States
Pontificia Universidade Católica do Rio de Janeiro (PUC-Rio), Rio de Janeiro, Brazil, associated to 2
University of Chinese Academy of Sciences, Beijing, China, associated to 3
School of Physics and Technology, Wuhan University, Wuhan, China, associated to 3
Institute of Particle Physics, Central China Normal University, Wuhan, Hubei, China, associated to 3
Departamento de Física, Universidad Nacional de Colombia, Bogota, Colombia, associated to 8
Institut für Physik, Universität Rostock, Rostock, Germany, associated to 12
National Research Centre Kurchatov Institute, Moscow, Russia, associated to 32
Instituto de Fisica Corpuscular, Centro Mixto Universidad de Valencia - CSIC, Valencia, Spain, associated to 38
Van Swinderen Institute, University of Groningen, Groningen, The Netherlands, associated to 43

Universidade Federal do Triângulo Mineiro (UFTM), Uberaba-MG, Brazil
