PROBING THE Earliest Stage of Protostellar Evolution—Barnard 1-bN and Barnard 1-bS

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ABSTRACT

Two submm/mm sources in the Barnard 1b (B1-b) core, B1-bN and B1-bS, have been observed with the Submillimeter Array (SMA) and the Submillimeter Telescope (SMT). The 1.1 mm continuum map obtained with the SMA reveals that the two sources contain spatially compact components, suggesting that they harbor protostars. The N2D+ and N2H+ J = 3–2 maps were obtained by combining the SMA and SMT data. The N2D+ map clearly shows two peaks at the continuum positions. The N2H+ map also peaks at the continuum positions, but is more dominated by the spatially extended component. The N2D+/N2H+ ratio was estimated to be ∼0.2 at the positions of both B1-bN and B1-bS. The derived N2D+/N2H+ ratio is comparable to those of the prestellar cores in the late evolutionary stage and the class 0 protostars in the early evolutionary stage. Although B1-bN is bright in N2H+ and N2D+, this source was barely seen in H13CO+. This implies that the depletion of carbon-bearing molecules is significant in B1-bN. The chemical property suggests that B1-bN is in the earlier evolutionary stage as compared to B1-bS with the H13CO+ counterpart. The N2H+ and N2D+ lines show that the radial velocities of the two sources are different by ∼0.9 km s−1. However, the velocity pattern along the line through B1-bN and B1-bS suggests that these two sources were not formed out of a single rotating cloud. It is likely that the B1-b core consists of two velocity components, each of which harbors a very young source.

Key words: stars: formation – stars: individual (B1-bN, B1-bS)

Online-only material: color figures

1. INTRODUCTION

The object formed at the beginning of the star formation process, the first hydrostatic core, was theoretically predicted by Larson (1969). Since the first hydrostatic core is in the transient stage between starless and protostellar cores, it is a key object to understand how stars are formed in molecular clouds. In the spherical symmetric case, first hydrostatic cores are extremely short lived, making the expected number of them only a few hundred of Class 0 protostars (Oomukai 2007; Masunaga & Inutsuka 2000). On the other hand, recent numerical simulations have shown that the rotating first cores can survive longer (Saigo et al. 2008; Saigo & Tomisaka 2006), providing more opportunities to observe them. Lately, some candidates for first cores have been reported (e.g., Chen et al. 2010, 2012; Enoch et al. 2010; Pineda et al. 2011); they have extremely low luminosity of <0.1 L⊙ and are not detectable at wavelengths shorter than ∼24 μm. However, the nature of these newly discovered sources has not been well understood yet.

The chemical properties are also useful tools to characterize the evolutionary stage of prestellar and protostellar cores. In the cold (∼10 K) and dense (∼104 cm−3) environment, carbon-bearing species such as CO and CS condense onto dust grains (Caselli et al. 2002; Tafalla et al. 2002). In contrast, nitrogen-bearing species remain in the gas phase longer (Bergin & Langer 1997). It is known that N2H+ and N2D+ lines are considered to be good tracers of cold dense gas because their abundances increase considerably if the CO is depleted. Observations of prestellar cores show that the spatial distributions of N2H+ and N2D+ well agree with that of dust continuum emission. In addition, in the cold and dense environment, the abundance of deuterium isotopologues is known to be enhanced significantly as compared to the elemental D/H ratio of ∼1.5 × 10−5 (Oliveira 2003). The N2D+/N2H+ ratio derived in a sample of prestellar cores ranges from 0.05 to 0.4 (Crapsi et al. 2005; Gerin et al. 2001). There is a tight correlation between N2D+/N2H+ ratio and the CO depletion factor (Crapsi et al. 2005). Once a protostar is formed in the core, the protostar warms up its environment and starts evaporating CO into the gas phase. As a result, the N2D+/N2H+ ratio drops as the protostar evolves (Emprechtinger et al. 2009). These results suggest that the N2D+/N2H+ ratio can be used as an indicator of evolution in both prestellar and protostellar cores. Especially protostellar cores with high N2D+/N2H+ ratio are particularly important because they are considered to be in a stage shortly after the formation of the central source.

In this study, we present the N2H+ J = 3–2 and N2D+ J = 3–2 observations with the Submillimeter Array (SMA) and the Submillimeter Telescope (SMT) toward two sources embedded in the Barnard 1 (B1) dark cloud (d = 250 ± 50 pc; Černis & Stražys 2003; Hirota et al. 2008), which is part of the Perseus molecular cloud complex. These two sources, labeled as Barnard 1-bN (B1-bN) and Barnard 1-bS (B1-bS) by Hirano et al. (1999), have no counterparts in the Spitzer images taken by the “From Molecular Cores to Planet Forming Disks” (c2d; Evans et al. 2003) survey in all four bands of IRAC, and 24 μm, 70 μm bands of MIPS (Jørgensen et al. 2004; Rebull et al. 2007). Both sources reveal the similar spectral energy distributions (SEDs), which show extremely low dust temperatures of <20 K (N. Hirano & F.-C. Liu, in preparation). Recently, Pezzuto et al. (2012) detected the far-infrared emission from B1-bN and B1-bS using the Herschel Space Observatory. On the basis of the SEDs, they proposed that B1-bN and B1-bS are first core candidates.

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2. OBSERVATIONS

2.1. Millimeter-interferometer Observations—SMA

SMA observations of the $N_2H^+ J = 3–2$ line were made on 2008 September 3. The array was in the subcompact configuration with seven antennas. The primary beam of the 6 m diameter antennas at 279 GHz was measured to be $\sim 45''$. The phase tracking center of the $N_2H^+ J = 3–2$ observations was at $\alpha(2000) = 3^h33^m21^s4$ and $\delta(2000) = 31^o07'35''3$. The spectral resolution mode was set to 101.6 kHz (1024 channels per 104 MHz chunk) for the $N_2H^+ J = 3–2$. The corresponding velocity resolution is 0.11 km s$^{-1}$. A pair of nearby quasars, 3C84 and 3C111, were used for gain calibration. Uranus was used for bandpass and absolute flux calibrations.

Observations of the $N_2D^+ J = 3–2$ line were carried out on 2007 September 10 and 11. The array was in the subcompact configuration. The primary-beam size of the antenna at 230 GHz was measured to be $\sim 54''$. The phase tracking center was at the same position as the $N_2H^+$ observations. The frequency resolution was set to be 101.6 kHz, which corresponds to a velocity resolution of 0.13 km s$^{-1}$. The $N_2D^+$ J = 3–2 line was observed simultaneously with the CO J = 2–1, $^{13}$CO J = 2–1, C$^{18}$O J = 2–1, and 230 GHz continuum, which will be presented in the separate paper by N. Hirano & F.-C. Liu (in preparation).

The visibility data were calibrated and edited using the MIRIAD software package. The calibrated visibility data were Fourier transformed and CLEANed with MIRIAD (Sault et al. 1995), using natural weighting. The synthesized beam sizes were 4′′53′′ × 3′′21′′ with a position angle (P.A.) of 44′′0 for the $N_2H^+ J = 3–2$ and 6′′08′′ × 3′′17′′ with a P.A. of 49′′2 for the $N_2D^+ J = 3–2$.

The 1.1 mm continuum data were obtained by averaging the line-free channels of the $N_2H^+ J = 3–2$ observations. The 1.1 mm continuum map had a resolution of 4′′53′′ × 3′′21′′ and an rms noise level of $\sim 0.35$ mJy beam$^{-1}$.

2.2. Single-dish Observations—SMT

Observations of the $N_2H^+ J = 3–2$ and $N_2D^+ J = 3–2$ lines toward the B1-b region were undertaken with the Arizona Radio Observatory SMT in 2008 November and December. In both lines, we made a 9 × 9 point map with a grid spacing of 10′. The map covers an area of 1.5′ × 1.5′ in right ascension and declination centered at $\alpha(2000) = 3^h33^m56^s$ and $\delta(2000) = 31^o09'34''8$. We used the 1.3 mm receiver system equipped with sideband separating mixers. The main-beam efficiency of the telescope was $\sim 0.74$. The half-power beam widths (HPBWs) were 27′′ and 33′′ in the $N_2H^+ J = 3–2$ and $N_2D^+ J = 3–2$ observations, respectively. For the $N_2H^+ J = 3–2$ observations, we used the chirp transform spectrometer, which provided a spectral resolution of 0.04 MHz. On the other hand, the $N_2D^+ J = 3–2$ spectra were obtained using the filter bank spectrometers, the spectral resolution of which was 0.25 MHz.

2.3. Combined Interferometric and Single-dish Observations

We combined the SMT data with the SMA data in order to fill the short-spacing information that was not sampled by the interferometer. We followed the procedure described by Yen et al. (2011), Takakuwa et al. (2007), and Wilner & Welch (1994). The combination of the data and subsequent imaging were made using the MIRIAD software package.

The single-dish image cube was Fourier transformed into the visibility data. Then the visibility data made from the single-dish image cube and those observed with the SMA were inverse Fourier transformed simultaneously into the image plane with natural weighting. The SMA+SMT combined image cube was made with a velocity interval of 0.3 km s$^{-1}$. The synthesized beam sizes of the combined maps were 4′′21′′ × 3′′83′′ with a P.A. of 49′′81′′ for the $N_2H^+ J = 3–2$ and 6′′40′′ × 4′′05′′ with a P.A. of 45′′75′′ for the $N_2D^+ J = 3–2$. The rms noise level of the combined $N_2H^+$ data cube was 0.17 Jy beam$^{-1}$ and that of the $N_2D^+$ data cube was 0.13 Jy beam$^{-1}$.

3. RESULTS

3.1. Dust Continuum Emission

Figure 1 shows the 1.1 mm continuum image of the B1-b region after the primary-beam correction. Two sources, B1-bN and B1-bS, separated by $\sim 20''$ are clearly seen. The peak flux densities, total integrated fluxes, and peak positions of B1-bN and B1-bS were determined by two-dimensional Gaussian fitting. The peak flux density at B1-bN is measured to be 0.28 ± 0.03 Jy beam$^{-1}$ and that at B1-bS is 0.45 ± 0.04 Jy beam$^{-1}$.

The beam deconvolved size of B1-bN is $2''2 \times 2''0$ (550 AU × 500 AU), and that of B1-bS is $2''0 \times 1''2$ (500 AU × 300 AU). The total fluxes of B1-bN and B1-bS are 0.37 Jy and 0.52 Jy, respectively. The positions of these compact sources are $\alpha(2000) = 3^h33^m21^s4$, $\delta(2000) = 31^o07'26''4$, and $\alpha(2000) = 3^h33^m21^s19$, $\delta(2000) = 31^o07'43''6$, respectively. The primary beam of the SMA also included the mid-infrared source located at $\sim 12''$ west of B1-bS. Although this source was detected in all four bands of IRAC and the 24 and 70 μm bands in MIPS (Jørgensen et al. 2006), there is no counterpart of this source in our 1.1 mm map. The 3σ upper limit of this source is $\sim 10$ mJy beam$^{-1}$.

We estimated the total mass of gas and dust in B1-bN and B1-bS by assuming the continuum emission to be optically thin at 1.1 mm, and a gas-to-dust ratio to be 100. The masses are estimated as follows:

$$M = \frac{F_s D^2}{\kappa_d B_s(T_{dust})},$$

where $F_s$ is the flux density obtained from continuum emission, $D$ is the source distance, and $\kappa_d$ is the dust mass opacity derived.
from $\kappa_0 = k_0 (v/v_0)^\beta$, where $k_0$ is calculated to be 0.01 cm$^2$ g$^{-1}$ at $k_0 = 231$ GHz (Ossenkopf & Henning 1994). We adopted the $\beta$ and $T_{\text{dust}}$ from the fits of SEDs in N. Hirano & F.-C. Liu (in preparation). The value of $\beta$ at B1-bN is 1.8 ± 0.4 and that at B1-bS is 1.3 ± 0.2. The values of $T_{\text{dust}}$ at B1-bN and B1-bS are 15.6 ± 2.2 K and 18.6 ± 1.6 K, respectively. The total masses of B1-bN and B1-bS are calculated to be 0.31 ± 0.20 $M_{\odot}$ and 0.34 ± 0.14 $M_{\odot}$, respectively.

The $H_2$ column densities at B1-bN and B1-bS were also estimated using the same assumptions of $\beta$, $T_{\text{dust}}$, and $k_0$. The resultant $N(H_2)$ at B1-bN and B1-bS are $(2.67 \pm 0.31) \times 10^{22}$ cm$^{-2}$ and $(4.84 \pm 0.72) \times 10^{22}$ cm$^{-2}$, respectively.

### 3.2. Molecular Line Maps

Figure 2 shows the integrated intensity maps of the $N_2H^+$ $J = 3–2$ and $N_2D^+$ $J = 3–2$ lines. The panels at the top of the figure are the maps obtained with only the SMA data, while the ones at the bottom include the contribution of the SMT data. The lines are integrated over the velocity ranges from $V_{\text{LSR}} = 5.4–8.7$ km s$^{-1}$ for the $N_2H^+$ $J = 3–2$ and from $V_{\text{LSR}} = 5.7–7.8$ km s$^{-1}$ for the $N_2D^+$ $J = 3–2$. The velocities of $N_2H^+$ $J = 3–2$ and $N_2D^+$ $J = 3–2$ are given with respect to the reference rest frequencies of 279,51169 GHz and 231,32990 GHz, respectively. The peaks of $N_2H^+$ $J = 3–2$ and $N_2D^+$ $J = 3–2$ emission agree well with those of 1.1 mm continuum emission in both SMA and SMA+SMT maps. In contrast, no significant $N_2H^+$ and $N_2D^+$ emission was detected at the position of the Spitzer source. From the comparisons between the line profiles of the SMA data and that of the combined SMA+SMT data, the missing fluxes of the SMA observations were estimated to be ~80% and ~67% for the $N_2H^+$ $J = 3–2$ and $N_2D^+$ $J = 3–2$, respectively. The larger missing flux in $N_2H^+$ implies that the $N_2H^+$ $J = 3–2$ emission is more spatially extended than the $N_2D^+$ $J = 3–2$ emission.

Figure 3 shows the velocity channel maps of the combined SMA+SMT $N_2H^+$ $J = 3–2$ data. In the low velocity range ($V_{\text{LSR}} = 6.0–6.6$ km s$^{-1}$), the $N_2H^+$ line emission is detected around both B1-bN and B1-bS regions. In this velocity range, the $N_2H^+$ emission reveals a centrally peaked distribution at B1-bS (especially at $V_{\text{LSR}} \sim 6.3$ km s$^{-1}$), while it is more spatially extended around B1-bN. In the high velocity range ($V_{\text{LSR}} = 6.9–7.8$ km s$^{-1}$), the emission is centrally peaked at B1-bN, while it is faint and spatially extended at B1-bS.

The velocity channel maps of the combined SMA+SMT $N_2D^+$ data (Figure 4) show the same trend as the $N_2H^+$ maps; the compact components at the positions of B1-bN and B1-bS appear in the velocity ranges of 6.9–7.8 km s$^{-1}$ and 6.0–6.6 km s$^{-1}$, respectively. As in the case of the $N_2H^+$, the spatially extended $N_2D^+$ emission around B1-bN is also seen in the low velocity range.

### 3.3. Kinematic Structure

In order to examine the velocity structure of B1-b, we made the position–velocity (P–V) diagrams of the $N_2H^+$ and $N_2D^+$ along the north–south cut through B1-bN and B1-bS (Figure 5). The upper panels of this figure are the diagrams obtained with the SMA data alone, while the ones in the lower panels include the contribution of the SMT data. The upper panels clearly show two velocity components; one is the 7.2 km s$^{-1}$ component at B1-bN and the other is the 6.3 km s$^{-1}$ component at B1-bS. At the position of B1-bN, the $N_2D^+$ emission has a secondary component at $V_{\text{LSR}} \sim 6.3$ km s$^{-1}$. This component
Figure 3. Velocity channel maps of the combined SMA+SMT N$_2$H$^+$ $J = 3$–$2$ data in B1-bN and B1-bS. The central velocity at each channel is shown at the top left corner. The small black oval at lower left in each panel indicates the synthesized beam size. The crosses are the peaks of 1.1 mm continuum emission and the peak of the Spitzer source. Contour levels start from 3σ by steps of 3σ for all figures, where 1σ is 0.17 Jy beam$^{-1}$.

is connected to the 6.3 km s$^{-1}$ component at B1-bS in the SMA+SMT combined map shown in the lower panels. The SMA+SMT combined map of N$_2$H$^+$ shows the faint 7.2 km s$^{-1}$ component at the position of B1-bS, however, this component is not seen in the N$_2$D$^+$ map.

4. ANALYSIS

4.1. N$_2$H$^+$ and N$_2$D$^+$ Hyperfine Fitting

Since N$_2$H$^+$ $J = 3$–$2$ and N$_2$D$^+$ $J = 3$–$2$ rotational transitions contain numerous hyperfine components, line parameters such as central velocity ($V_c$), line width ($\Delta v$), optical depth ($\tau$), and excitation temperature ($T_{\text{ex}}$) need to be determined by hyperfine fitting. We assume all the hyperfine components to have the same excitation temperature and line width, with the relative velocity of each component fixed to the laboratory value. Since B1-bN has two velocity components (Figure 5), it was difficult to determine all the parameters from fitting. Therefore, we fixed the excitation temperatures to be 7.6 K for N$_2$H$^+$ and 6.7 K for N$_2$D$^+$, which were determined by Emprechtinger et al. (2009). Since these excitation temperatures were derived from the N$_2$H$^+$ $J = 1$–$0$ and N$_2$D$^+$ $J = 1$–$0$ single-dish observations with HPBW of $\sim$27$''$ and 32$''$, respectively, the values are considered to be the average values of the B1-b region. We adopted the N$_2$H$^+$ $J = 3$–$2$ hyperfine components listed in
Figure 4. Velocity channel maps of the combined SMA+SMT N2D+ J = 3–2 data in B1-bN and B1-bS. The central velocity at each channel is shown at the top left corner. The small black oval at lower right in each panel indicates the synthesized beam size. The crosses are the peaks of 1.1 mm continuum emission and the peak of the Spitzer source. Contour levels start from 3σ by steps of 3σ for all figures, where 1σ is 0.13 Jy beam−1.

Caselli et al. (2002) and those of the N2D+ J = 3–2 listed in Gerin et al. (2001). We used the combined SMA+SMT data to perform the fitting. The N2H+ data were convolved in order to match the beam size of the N2D+ data (6′.40 × 4′.05 with a P.A. of 45°75). The primary beam of the SMA antennas was also corrected.

Figure 6 shows the combined SMA+SMT N2H+ J = 3–2 and N2D+ J = 3–2 line profiles and the results of hyperfine fittings at the positions of B1-bN and B1-bS. The N2D+ line profiles were fitted well using the optically thin assumption, while the N2H+ lines were obviously optically thick. The spectra at B1-bN have two velocity components; one is at \( V_{\text{LSR}} \approx 7.2 \text{ km s}^{-1} \), which is spatially compact, and the other is the spatially extended component at \( V_{\text{LSR}} \approx 6.2 \text{ km s}^{-1} \). The line profiles at B1-bS are mostly dominated by the \( V_{\text{LSR}} = 6.3 \text{ km s}^{-1} \) component. The determined parameters are listed in Tables 1 and 2. In these tables, the main components refer to the dominant components at B1-bN and B1-bS, which correspond to the velocity components centered at \( \sim 7.2 \text{ km s}^{-1} \) and \( 6.3 \text{ km s}^{-1} \), respectively. Since the line profile of N2H+ J = 3–2 at the position of B1-bS shows the secondary component at \( \sim 7.2 \text{ km s}^{-1} \), we included this component to the fitting.

It should be noted that the N2H+ line becomes optically thin if the higher excitation temperature is assumed; the optical depth of the primary component of B1-bN becomes less than 1.0 for the \( T_{\text{ex}} > 8.8 \text{ K} \), and that of B1-bS becomes <1 for the \( T_{\text{ex}} > 10.5 \text{ K} \). On the other hand, the optical depth of the N2H+ line becomes the largest (\( \sim 10 \) for B1-bN and \( \sim 18 \) for B1-bS) if we adopt the lower limit of the excitation temperature derived from the fitting (\( \sim 6.4 \text{ K} \) for B1-bN and \( \sim 7.2 \text{ K} \) for B1-bS).

4.2. N2H+ and N2D+ Column Densities

We estimated the column densities of N2H+ and N2D+ using the method described by Caselli et al. (2002). The results including column densities of N2H+ and N2D+ and the \( N(N_2D^+)/N(N_2H^+) \) ratios at B1-bN and B1-bS are summarized in Table 3. The column densities of N2H+ at the positions of B1-bN and B1-bS were calculated to be \( \sim 1.1 \times 10^{13} \text{ cm}^{-2} \) and \( 1.8 \times 10^{13} \text{ cm}^{-2} \), respectively, while column densities of N2D+ at B1-bN and B1-bS are \( \sim 2.9 \times 10^{12} \text{ cm}^{-2} \) and \( 2.4 \times 10^{12} \text{ cm}^{-2} \), respectively. The corresponding \( N(N_2D^+)/N(N_2H^+) \) ratios were \( \sim 0.27 \) at B1-bN and \( \sim 0.13 \) at B1-bS. These values are consistent with the \( N_2D^+/N_2H^+ \) ratio of \( \sim 0.18 \) in the entire B1-b region obtained by Emprechtinger et al. (2009). The N2H+ column density becomes lower if we adopt higher excitation temperatures (a factor of 6–9 if \( T_{\text{ex}} = 20 \text{ K} \)). However, the N2D+/N2H+ ratio is almost constant at \( \sim 0.3–0.4 \), because both N2H+ and N2D+ lines are optically thin at \( T_{\text{ex}} \gtrsim 10 \text{ K} \).
Figure 5. P–V diagrams of (a) SMA N$_2$H$^+$ $J = 3–2$, (b) SMA N$_2$D$^+$ $J = 3–2$, (c) combined SMA+SMT N$_2$H$^+$ $J = 3–2$, and (d) combined SMA+SMT N$_2$D$^+$ $J = 3–2$ emission along the N–S cut through B1-bN and B1-bS. The horizontal lines indicate the positions of B1-bN (orange) and B1-bS (purple). Contour levels are from 2σ in steps of 2σ, where 1σ are 0.17 and 0.13 Jy beam$^{-1}$ in the N$_2$H$^+$ and N$_2$D$^+$ data, respectively. (A color version of this figure is available in the online journal.)

Table 1

| Core     | Main Component | Second Component |
|----------|----------------|------------------|
|          | $\tau_{tot}$  | $V_c$ (km s$^{-1}$) | $\Delta v$ (km s$^{-1}$) | $\tau_{tot}$  | $V_c$ (km s$^{-1}$) | $\Delta v$ (km s$^{-1}$) |
| B1-bN    | 2.00 ± 0.03   | 7.19 ± 0.02      | 0.83 ± 0.05                 | 0.82 ± 0.02   | 6.18 ± 0.02      | 0.50 ± 0.05                 |
| B1-bS    | 8.00 ± 0.06   | 6.27 ± 0.02      | 0.35 ± 0.03                 | 0.10 ± 0.08$^a$| 7.25 ± 0.01      | 0.21 ± 0.02                 |

Note. $^a$ Though the second component is not clear at B1-bS, two components were used to fit the spectrum of N$_2$H$^+$ $J = 3–2$ at B1-bS because it seems to be a more spatially extended component centered at ~6.2 km s$^{-1}$.

On the other hand, if we adopt the lower limit of the excitation temperature, the column density of N$_2$H$^+$ increases by a factor of ~3–4. Since the N$_2$D$^+$ line is assumed to be optically thin, the N$_2$D$^+$/N$_2$H$^+$ ratio decreases to ~0.06.

The fractional abundances of N$_2$H$^+$ at B1-bN and B1-bS were calculated to be $3.97 \pm 0.62 \times 10^{-10}$ and $3.71 \pm 0.63 \times 10^{-10}$, respectively, using the N$_2$H$^+$ column densities and H$_2$ column densities derived from the continuum data. Jørgensen
Figure 6. Line profiles of (a) N$_2$H$^+$ $J = 3–2$ and (b) N$_2$D$^+$ $J = 3–2$ at the position of B1-bN, (c) N$_2$H$^+$ $J = 3–2$, and (d) N$_2$D$^+$ $J = 3–2$ at that of B1-bS. The dashed curves (orange) indicate individual velocity components, while the solid curves (brown) show the overall fitting results. (A color version of this figure is available in the online journal.)

Table 2

Results of N$_2$D$^+$ $J = 3–2$ Hyperfine Fitting

| Core   | $I$ (a) (K) | $V_c$ (km s$^{-1}$) | $\Delta v$ (km s$^{-1}$) | $I$ (K) | $V_c$ (km s$^{-1}$) | $\Delta v$ (km s$^{-1}$) |
|--------|-------------|---------------------|--------------------------|--------|---------------------|--------------------------|
| B1-bN  | 1.30 ± 0.06 | 7.33 ± 0.02         | 0.94 ± 0.05              | 0.87 ± 0.08 | 6.28 ± 0.03         | 0.69 ± 0.04              |
| B1-bS  | 1.80 ± 0.03 | 6.43 ± 0.03         | 0.58 ± 0.05              |        |                     |                          |

Note. a The brightness temperature of each hyperfine component is $I \times$ (relative intensity of each component).

Table 3

N$_2$H$^+$, N$_2$D$^+$, and H$_2$ Column Densities

| Core   | N(N$_2$H$^+$) (cm$^{-2}$) | N(N$_2$D$^+$) (cm$^{-2}$) | N(N$_2$D$^+$)/N(N$_2$H$^+$) | N(H$_2$) (cm$^{-2}$) | N(N$_2$H$^+$)/N(H$_2$) |
|--------|--------------------------|--------------------------|-----------------------------|---------------------|--------------------------|
| B1-bN  | (1.06 ± 0.11) x 10$^{13}$ | (2.85 ± 0.09) x 10$^{12}$ | 0.27 ± 0.11                 | (2.67 ± 0.31) x 10$^{22}$ | (3.97 ± 0.62) x 10$^{-10}$ |
| B1-bS  | (1.80 ± 0.15) x 10$^{13}$ | (2.40 ± 0.09) x 10$^{12}$ | 0.13 ± 0.09                 | (4.84 ± 0.72) x 10$^{22}$ | (3.71 ± 0.63) x 10$^{-10}$ |

et al. (2004) derived the abundance of N$_2$H$^+$ in prestellar cores and Class 0 sources to be $\sim$10$^{-8}$–10$^{-10}$. The N$_2$H$^+$ abundance in these two sources is comparable to those of the samples studied by Jørgensen et al. (2004).

5. DISCUSSIONS

5.1. Deuterium Fractionation

The N$_2$D$^+/N$_2$H$^+$ ratio of 0.13–0.27 derived in the previous section is four orders of magnitude larger than the D/H ratio of 10$^{-5}$ in the interstellar space (Oliveira 2003). It is known that the deuterium fractionation in N$_2$H$^+$ is related to the core temperature and CO depletion (e.g., Di Francesco et al. 2007). In the proton–deuteron exchange reaction, H$_3^+$ + HD $\rightarrow$ H$_2$D$^+$ + H$_2$, H$_2$D$^+$/H$_3^+$ tends to be increased at lower temperature (Millar et al. 1989). In addition, at densities of $\gtrsim$10$^5$ cm$^{-3}$, the depletion of carbon-bearing molecules onto dust grains further increases the H$_2$D$^+$/H$_3^+$ ratio (Roberts & Millar 2000). Since H$_2$D$^+$ is the main ingredient of deuterated molecules, the enhancement of H$_2$D$^+$ increases the abundance of deuterated molecules. Therefore, the N$_2$D$^+/N$_2$H$^+$ ratio becomes an indicator of the physical conditions and molecular depletion of dense and cold gas, and is useful to probe the chemical and dynamical evolution of cores, including starless cores and class 0 protostars.

Crapsi et al. (2005) observed 31 low-mass starless cores and derived the N$_2$D$^+/N$_2$H$^+$ ratio to be 0.03–0.44. They found a strong correlation between the N$_2$D$^+/N$_2$H$^+$ ratio and CO depletion, suggesting that the N$_2$D$^+/N$_2$H$^+$ ratio increases with the density of the core. The higher N$_2$D$^+/N$_2$H$^+$ values were seen in the cores of later evolutionary stage, i.e., those with centrally concentrated density distribution and high H$_2$ column densities. In the case of the Class 0 protostellar cores studied by Emprechtinger et al. (2009), the N$_2$D$^+/N$_2$H$^+$ ratio also ranges from 0.03 to 0.3. Their study showed that the N$_2$D$^+/N$_2$H$^+$ ratio in the prestellar cores significantly decreases as a function of dust temperature, suggesting that the ratio decreases...
as the central star evolves and heats the surrounding gas. These results suggest that the largest $N_2D^*/N_2H^+$ ratio is expected to be in prestellar cores that are close to forming protostars or protostellar cores in the earliest evolutionary stage. The $N_2D^*/N_2H^+$ ratios derived at B1-bN and B1-bS ($\sim 0.2$) are comparable to those of the evolved prestellar cores such as L1544 and the very young protostars such as L1448C, IRAS 03282, and HH211. It is notable that the column densities of $N_2D^+$ estimated here are the lower limits, because we assumed the $N_2D^+$ line to be optically thin. Therefore, the $N_2D^*/N_2H^+$ ratios derived here are also the lower limits.

The high D/H ratios in the B1-b region are also shown in other observations (Lis et al. 2002; Roueff et al. 2005; Marcelino et al. 2005). The high D/H ratios are consistent with the low temperatures derived from the SEDs of the continuum emission (Pezzuto et al. 2012; N. Hirano & F.-C. Liu, in preparation). Since both B1-bN and B1-bS contain compact continuum sources that are clearly detected by the interferometer, they are unlikely to be the prestellar cores. On the basis of the properties of B1-bN and B1-bS derived from the continuum SEDs such as low temperatures ($T_{\text{dust}} < 20$ K), low intrinsic luminosities ($L_{\text{int}} \sim 0.03 L_\odot$ for B1-bS and $\lesssim 0.004 L_\odot$ for B1-bN), and low-velocity CO outflows (N. Hirano & F.-C. Liu, in preparation), these two sources are considered to be the candidates for the first hydrostatic cores.

### 5.2. Kinematics

As shown in Section 3.3, B1-bN appears at $V_{\text{LSR}} = 7.2$ km s$^{-1}$ and B1-bS reveals at $V_{\text{LSR}} = 6.3$ km s$^{-1}$. In other words, two sources with a projected separation of $\sim 5000$ AU have a velocity difference of $\sim 0.9$ km s$^{-1}$. If this velocity difference attributes to the orbital motion of gravitationally bounded objects, the mass of each source is estimated to be $\sim 0.55 M_\odot$. This mass agrees with the mass derived from the continuum data within a factor of two. If the two sources originate from a single rotating cloud, the P–V diagram along the line connecting two sources (i.e., in the N–S direction) is expected to show a velocity gradient. However, the velocity pattern seen in the P–V maps including the short-spacing data (Figures 5(c) and (d)) does not show such a velocity gradient. Instead, these P–V maps reveal a C-shaped pattern with two velocity components along the line of sight of B1-bN. The spatially extended emission connects B1-bS and the secondary component of B1-bN, which has the same velocity as B1-bS. In addition, the $N_2H^+$ map shows an additional 7.2 km s$^{-1}$ component along the line of sight of B1-bN. Therefore, it is unlikely that the two sources were formed in a single rotating cloud through the fragmentation. The P–V maps including the short-spacing data suggest that there are two velocity components along the line of sight of the B1-b core, and that B1-bN belongs to the 7.2 km s$^{-1}$ cloud and B1-bS to the 6.3 km s$^{-1}$ cloud.

### 5.3. Comparison with $H^{13}CO^+ J = 1–0$ Emission

Figure 7 shows the P–V diagram of $H^{13}CO^+ J = 1–0$ emission along the N–S cut through B1-bN and B1-bS observed within the NRO 45 m telescope (N. Hirano & F.-C. Liu, in preparation). Two velocity components along the line of sight of B1-b are also seen in the $H^{13}CO^+$ map. It is notable that the relative intensities of these two velocity components in the $H^{13}CO^+$ map are considerably different from those of the $N_2H^+$ and $N_2D^+$ maps. As shown in the previous section, the dense gas around B1-bN traced by the $N_2H^+$ and $N_2D^+$ lines has a velocity of 7.2 km s$^{-1}$. In addition to this 7.2 km s$^{-1}$ component, there is a secondary component at 6.3 km s$^{-1}$. On the other hand, the $H^{13}CO^+$ line peaks at 6.3 km s$^{-1}$ and shows no significant enhancement at $V_{\text{LSR}} \sim 7.2$ km s$^{-1}$. At B1-bS, both $N_2H^+$ and $N_2D^+$ maps reveal the 6.3 km s$^{-1}$ component. Although the $N_2H^+$ line shows a weak secondary peak at $\sim 7$ km s$^{-1}$, there
is no counterpart of that velocity component in the N$_2$D$^+$ line. In the case of H$_3$CO$^+$ map, however, it shows double peaks at 6.3 km s$^{-1}$ and 6.9 km s$^{-1}$. The comparison shows that N$_2$H$^+$ and N$_2$D$^+$ emission is partially anti-correlated with the H$_3$CO$^+$ emission.

These results can be well explained by the different chemical properties between carbon-bearing molecules and nitrogen-bearing molecules. Since H$_3$CO$^+$ (HC$_3$O$^+$) is considered to be formed from $^{13}$CO (CO) through the ion–neutral reaction in the gas phase, the abundance of H$_3$CO$^+$ is expected to decrease if a significant fraction of CO molecules freeze onto the dust grains. In such a cold and dense region with CO depletion, the N$_2$H$^+$ and N$_2$D$^+$ abundance tends to be enhanced, since CO is the main reactant to destroy the N$_2$H$^+$ and N$_2$D$^+$. Therefore, the chemical properties of the dense gas surrounding B1-bN with a velocity of 7.2 km s$^{-1}$ can be explained by the scenario of CO depletion. In the case of B1-bS, the H$_3$CO$^+$ coexists with N$_2$H$^+$ and N$_2$D$^+$. This is probably because the gas in B1-bS is heated by the newly formed star. As a consequence, the CO in this region has partially come back to the gas phase. The presence of CO outflows around this source also supports this idea (N. Hirano & F.-C. Liu, in preparation). Since H$_{13}$CO$^+$ is made from 13CO (CO) through the ion–neutral reaction in the gas phase, the abundance of H$_{13}$CO$^+$ is bright in the N$_2$H$^+$ and N$_2$D$^+$, is not clearly seen in the H$_3$CO$^+$ J = 1–0. This suggests that the carbon-bearing molecules in the dense gas surrounding B1-bN are heavily depleted onto the dust.

5. The radial velocities of B1-bN and B1-bS are $v_{LSR} = 7.2$ km s$^{-1}$ and $v_{LSR} = 6.3$ km s$^{-1}$, respectively. The C-shaped velocity pattern seen in the P–V diagrams of N$_2$H$^+$ and N$_2$D$^+$ does not support the idea that two sources are formed by means of fragmentation of a single rotating cloud. It is likely that there are two velocity components at $v_{LSR} = 6.3$ km s$^{-1}$ and 7.2 km s$^{-1}$ along the line of sight of B1-b, each of which harbors a young protostar.

6. At the position of B1-bN, the 7.2 km s$^{-1}$ component, which is bright in N$_2$H$^+$ and N$_2$D$^+$, is not clearly seen in the H$_3$CO$^+$ J = 1–0. This suggests that the carbon-bearing molecules in the dense gas surrounding B1-bN are heavily depleted onto the dust.

7. The on the other hand, the 6.3 km s$^{-1}$ component at B1-bS is observed in N$_2$H$^+$, N$_2$D$^+$, and H$_{13}$CO$^+$ lines. The difference in chemical property suggests that B1-bN is in the earlier evolutionary stage as compared to B1-bS.

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