Comparison of computational chemistry methods for the discovery of quinone-based electroactive compounds for energy storage

Qi Zhang1,2,3, Abhishek Khetan1,2 & Süleyman Er1,2

High-throughput computational screening (HTCS) is a powerful approach for the rational and time-efficient design of electroactive compounds. The effectiveness of HTCS is dependent on accuracy and speed at which the performance descriptors can be estimated for possibly millions of candidate compounds. Here, a systematic evaluation of computational methods, including force field (FF), semi-empirical quantum mechanics (SEQM), density functional based tight binding (DFTB), and density functional theory (DFT), is performed on the basis of their accuracy in predicting the redox potentials of redox-active organic compounds. Geometry optimizations at low-level theories followed by single point energy (SPE) DFT calculations that include an implicit solvation model are found to offer equipollent accuracy as the high-level DFT methods, albeit at significantly lower computational costs. Effects of implicit solvation on molecular geometries and SPEs, and their overall effects on the prediction accuracy of redox potentials are analyzed in view of computational cost versus prediction accuracy, which outlines the best choice of methods corresponding to a desired level of accuracy. The modular computational approach is applicable for accelerating the virtual studies on functional quinones and the respective discovery of candidate compounds for energy storage.

Commercial utilization of intermittent renewable energy sources, such as solar and wind, requires large-scale, low-cost, and durable energy storage technologies to balance the mismatch between energy supply and demand. Redox flow batteries (RFBs) are recognized as prime candidates for large-scale and variable-term storage of electrical energy1–3. RFBs have external storage tanks that store the liquid-phase redox-active electrolyte materials separated from the electrochemical reaction cells. This unique design feature is advantageous as it decouples the battery's power and energy density scaling, while also facilitating easier maintenance and recycling2.

Conventional RFBs are operated using metal-based electrolyte materials, such as vanadium, iron, zinc, lead and chromium4. They, however, face technical challenges of ion crossover through the membranes and sluggish reaction kinetics5. Additionally, the cost of active materials and the risks associated with metal toxicity have been hindering widespread deployment of metal-based RFBs6,7. RFBs employing organic redox-active materials offer a promising alternative to metal-based electrolytes, as they can be sourced from earth-abundant elements and modified further structurally to tune their key battery-relevant properties1,2. The emerging classes of organic redox-active RFB compounds consist of quinones1, viologens8,9, alloxazines10,11, phenazines12,13, and nitroxide radicals14. Quinones are ubiquitous in nature15, and with their fast redox kinetics16,17 and tunable properties owing to their chemical diversity1,18, they are increasingly being utilized as electroactive materials in advanced RFB technologies. In recent years, an increasing effort has been made to develop aqueous RFBs (ARFBs) that use quinones as electroactive materials, including the functionalized forms of benzoquinones19,20, naphthoquinones17,19, and anthraquinones9,21. Research has shown that these molecules undergo a coupled two-electron two-proton redox reaction in aqueous media22. However, these molecules offer low energy densities in practical ARFBs as they are not very soluble and their half-cell redox potentials are not close to 0 V versus SHE.

1DIFFER—Dutch Institute for Fundamental Energy Research, De Zaale 20, 5612 AJ Eindhoven, The Netherlands. 2CCER—Center for Computational Energy Research, De Zaale 20, 5612 AJ Eindhoven, The Netherlands. 3Department of Applied Physics, Eindhoven University of Technology, 5600 MB Eindhoven, The Netherlands. *email: s.er@differ.nl
which is desired for ARFB anolytes\textsuperscript{19,23,24}. Therefore, a major challenge for organic ARFBs is to tune the properties of the electroactive compounds to meet the practical requirements of high power and high energy density batteries. To develop an ARFB with a large cell voltage, maximizing the redox potential window of quinone-based compounds is essential. Recent experimental and computational studies show that the redox potential of organic ARFBs are significantly influenced by functionalizing them with electron-withdrawing/donating groups\textsuperscript{1,2}. Assary et al.\textsuperscript{25} and Aspuru-Guzik et al.\textsuperscript{26,27} used HTCS methods for creating virtual libraries of candidate electroactive compounds populated with the functionalized compounds of quinones and predicting their redox properties. These studies utilized robust quantum chemical calculations to estimate the thermodynamic properties of compounds and identify the most promising candidates. Thus, they demonstrated the usefulness of hierarchical HTCS methods in accelerating the property predictions of redox-active molecules. Using quantum chemical calculations to predict the redox properties is, however, a practically challenging task. The approach is particularly not well-suited for HTCS studies on a huge space of conceivable molecules. Therefore, there is a need to determine the trade-offs between the prediction accuracy and the computational cost. While there has been a significant increase in the number of HTCS efforts for RFBs\textsuperscript{25–32}, to the best of our knowledge, an analysis of the factors that affect prediction accuracy, such as the level of theory for optimization of molecular geometry, inclusion/exclusion of solvation effects, and the level of theory for the calculation of chemical descriptors, are not available in the current literature.

Here, we systematically evaluate the performance of different computational methods, including DFT, DFTB\textsuperscript{29}, and SEQM\textsuperscript{30,31}. We compare them based on their accuracy in predicting the experimentally measured redox potentials of quinones obtained from four different sources. In addition, we make first-order comparisons of the computational cost of these approaches to offer the best approach for the prediction of redox potentials. Our results provide new insights on the factors that influence the prediction performance of computational methods. The findings are expected to be useful for both customary and HTCS efforts that are aimed at studying redox-active molecules for RFBs, and also beyond ARFB compounds that take part in bio- and electro-chemical conversion reactions.

Computational workflow. We developed a systematic workflow to make generalizable and consistent comparisons between the different computational approaches. As shown in Fig. 1, the starting point in the workflow for any molecule is its SMILES representation\textsuperscript{33}, which is a widely used form of graph representation and can easily be generated for any given molecule. The SMILES representation is at first converted to a two-dimensional (2D) geometrical representation using a SMILES interpreter. Next, (1) The 2D representation is converted to a three-dimensional (3D) geometry by applying the geometry optimization (OPT) scheme of OPLS\textsubscript{3e} FF\textsuperscript{36} to identify the lowest energy 3D conformer. As shown in Fig. 1b, the FF level geometry is the starting point for all approaches. (2) The 3D geometry is further optimized in the gas-phase at three different levels of theory, namely: SEQM, DFTB, and DFT. For SEQM and DFT, geometry optimizations have also been performed in an implicit aqueous-phase, but for simplicity they are not shown in Fig. 1. This step yields different 3D geometries and the corresponding SPE of the compounds. (3) Next, SPEs of all the different 3D geometries are calculated using different DFT functionals. This step yields energy values that are directly comparable but are obtained from geometry optimizations that have been performed at four different levels of theory. (4) Lastly, for the geometries obtained from the optimizations in the gas-phase, the SPEs are recalculated, this time by including the effect of the aqueous medium (SOL) implicitly by using the Poisson–Boltzmann solvation model (PBF)\textsuperscript{37,38}.

Results and discussions
Comparison of DFT methods. Since DFT is the highest level of theory considered in the current study, we begin with a discussion of the performance of the various DFT functionals, also with an aim to use them as performance benchmarks for low-level methods. Initially, we briefly discuss the performance of $\Delta U_{\text{rxn}}$ and $\Delta G_{\text{rxn}}$ as chemical descriptors to predict redox potentials. For this purpose, DFT energy calculations using the PBE functional are performed, first for optimizing geometries in gas-phase and then for calculating single point energies in an implicit aqueous-phase. The calibration performances of $\Delta U_{\text{rxn}}$ (RMSE = 0.049 V, $R^2$ = 0.978) and $\Delta G_{\text{rxn}}$ (RMSE = 0.048 V, $R^2$ = 0.979) are very similar, as shown in Supplementary Fig. S1. Inclusion of ZPE in $\Delta U_{\text{rxn}}$, as well as entropic effects in $\Delta G_{\text{rxn}}$, is only marginally better than using $\Delta E_{\text{rxn}}$ (RMSE = 0.051 V, $R^2$ = 0.977). Therefore, we consider that the effects of including these terms are not significant enough from HTCS perspective. Accordingly, all the following discussions in this work consider $\Delta E_{\text{rxn}}$ as the descriptor.

First, we discuss linear calibrations of three representative DFT functionals of PBE (Fig. 2a), B3LYP (Fig. 2b), and M08-HX (Fig. 2c). The results shown in Fig. 2 have been obtained by using three kinds of DFT calculated reaction energies $\Delta E_{\text{rxn}} = \Delta E_{\text{DFT}}$ against the experimentally measured redox potentials ($E_{\text{exp}}^o$) as follows: (1) OPT in gas-phase without calculation of SPE in SOL, (2) OPT in gas-phase and following SPE in SOL, and (3) both OPT and SPE in SOL.

On comparing RMSE and $R^2$ data, the following observations are made:

(1) When using $\Delta E_{\text{DFT}}$ from only gas-phase optimized geometry and SPE, PBE is the least accurate functional at the GGA level with RMSE = 0.072 V, $R^2$ = 0.954. Nevertheless, the results show that at any DFT level, it is possible to predict $E_{\text{exp}}^o$ for quinone-based molecules within a range of common experimental errors ($\approx 0.1$ V).
Upon the inclusion of solvation effects using an implicit model on the gas-phase geometries, all the three functionals show a decrease in their RMSE values. The percentage decrease in error is highest for PBE (30%) and lowest for M08-HX (23%).

Remarkably for all the three considered functionals, full geometry optimizations and energy calculations in an implicit solvation model yield slightly worse results than their counterparts in which geometries are optimized in gas-phase. The RMSEs increase by 0.002–0.004 V, indicating that there is no real added value of performing geometry optimizations with implicit solvation, not to mention that they are also computationally more demanding.

Based on the findings above, we evaluate the performances of eight other DFT functionals without considering geometry optimizations in an implicit solution-phase. In Fig. 2d and Supplementary Fig. S2, a summary of the performance of all the DFT functionals considered in this work is presented using bar plots for RMSE and $R^2$, respectively. When compared under the same set of conditions, all functionals, with the exception of LDA, show a similar performance. The PBE0/PBE0-D3, HSE06 and M08-HX functionals show a highly similar performance when using implicit solvation on gas-phase optimized geometries, which are followed by the other hybrid functionals.

Figure 1. (a) Schematic showing the chain of operations for geometry optimization (OPT) and single point energy (SPE) calculations at different levels of theory (b) A graphical summary of the various levels of approximations used to estimate $\Delta E_{\text{DFT}}^\text{rxn}$. The hollow black arrows with symbol $\Delta$ represent the difference between gas-phase reaction energies of the optimized geometries obtained at different levels of theory. The solid black arrows with symbol $\Delta$ represent the difference between energies computed using DFT on a fixed geometry obtained from a lower-level theory and energies computed at that given level of theory. The dotted gray arrows represent the solvation effect from gas-phase SPE to solution-phase SPE, when the implicit solvation model is considered. For both (a,b), the text boxes with no background color represent geometry optimizations; the boxes with colored background represent SPE calculations; the boxes with water bubbles in the background represent solution-phase SPE calculations using the implicit aqueous solvent model.

$\Delta$ [RMSE] in predictions of $E_{\text{exp}}^0$ values between:

$\Delta^\text{a}$ – gas phase energy and geometry at different levels of theory

$\Delta^\text{e}$ – energy computed at DFT level for a fixed geometry from a lower level theory

$\cdots$ $\cdots$ $\Delta^\text{a}$ – implicit solvation at DFT level
functionals of B3LYP/B3LYP-D3, and then by the PBE/PBE-D3 and PBE/PBE-D3 GGA functionals. The addition of D3-dispersion corrections makes hardly any difference on either of the hybrid or the GGA functional calculated results. For all further comparisons in this work, we choose, among the compared exchange–correlation functionals, PBE as the benchmark DFT functional as it offers the best compromise between prediction accuracy and computational cost. We note that, the different DFT functionals have been compared here purely on the basis of their performance in predicting the measured potentials. We also note that, functionals constructed with higher degrees of empiricism, such as the Minnesota density functionals, are aimed at producing better values for a chosen set of physically observable properties. In this regard, it is not surprising that the M08-HX performs best among the functionals, as it is heavily parametrized to show good performance for thermochemistry. However, it must also be kept in mind that such heavily parametrized functionals tend to produce less accurate electron densities than the ones with little to no empiricism in their design (e.g. the PBE functional).
Discussion
As shown in Fig. 2d, the DFT LDA data is not as good as the GGA and the hybrid method calculated data. Additionally, there is a positive impact on the prediction accuracies due to the inclusion of implicit solvation in calculations of $\Delta E_{\text{rxn}}$. The effect can be attributed to a better accounting of the $-\text{OH}$ groups' interactions with the surrounding aqueous environment in the hydroquinone products\(^\text{37}\). Surprisingly, optimizing geometries with implicit solvation slightly worsens the prediction accuracies. This can be attributed to multiple factors. First, it is possible that the PBF solvation model is not accurate enough to improve the gas-phase geometry. Secondly, there might be a serendipitous cancellation of errors when using the gas-phase geometry that is affected by the changes in the geometry due to the implicit solvation model in use. Additionally, $\Delta E_{\text{rxn}}$ is used as an approximation for $\Delta G_{\text{cal}}$ and accounting for the ignored pressure–volume and entropy terms from Eq. (6) might result in better prediction accuracies when optimizing the geometries in solution. In the work of Kim et al.\(^\text{32}\) it was shown that the reduction potentials of Anthraquinones in acidic aqueous solutions are strongly influenced by specific interactions with molecules in solvent environment. In aqueous solution, they found that using DFT (ωB97X-D/6-31G*) with implicit solvation (PCM(Bondi)) for geometry optimizations yields good results, except for the redox couples that have strong intramolecular hydrogen bond interactions. They evaluated a total of 19 Anthraquinones and reported a mean absolute deviation (MAD) of 0.194 V for three outliers that showed strong intramolecular hydrogen bond interactions. This value was more than five times the MAD value of the remaining 16 redox couples (0.037 V). Further, they showed that QM/MM calculations (with the TIP3P force field for explicit water molecules) alleviate the overestimation and lead to a more balanced treatment of solute–solvent interactions. Accordingly, using a QM/MM model, the correlation between theory and experiment data had a MAD of 0.033 V. In the current work, we performed a similar analysis on our calibration data set of 43 molecule pairs. We found that there are only 12 molecules, with IDs: 1, 2, 3, 4, 5, 6, 8, 9, 16, 35, 37 and 39 shown in Supplementary Table S1, without any possibility of strong intramolecular hydrogen bond interactions due to the neighboring positions found in the hydroquinone versions of the molecules. Surprisingly, we found that when using implicit aqueous solvation during geometry optimizations, the MADs for the predicted redox potentials were 0.039 V for the 12 molecules and 0.037 V for the remaining 31 molecules. We note that these MADs are very similar and the difference between the two groups is only in the third decimal digit. Therefore, we cannot confirm that the explanation provided by Kim et al. also applies to the methods used in this work. At the same time, it must be noted that Kim et al. used only Anthraquinones (3-ring molecules) for their analysis, whereas this work considers a wide variety of quinone molecules (from 1 to 3 rings), including those with the C=O groups at the 1,2 positions on the compounds. Further, Kim et al. employed the PCM (Bondi) implicit solvation model, which is different from the PBF model used in the current work. These differences, as well as the difference in the calibration data, make it hard to ascertain the exact origin of the disparities between this work and the work of Kim et al.

Another important aspect of the calibration of molecules that needs to be considered is the effect of ionization of sulfonic acid groups, as they are prone to dissociation in aqueous media. In the calibration set of 43 molecule pairs used in the current work, there are 18 molecules that contain $-\text{SO}_3\text{H}$ groups, with IDs: 7, 8, 9, 10, 11, 12, 13, 30, 31, 32, 33, 34, 35, 36, 37, 41, 42 and 43 as shown in Supplementary Table S1. In the framework of the best performing scheme, i.e. PBE, the calculated MAD for these 18 molecules is 0.047 V, which is approximately 13, 30, 31, 32, 33, 34, 35, 36, 37, 41, 42 and 43 as shown in Supplementary Table S1. In the work of Kim et al.\(^\text{32}\) it was shown that the reduction potentials of Anthraquinones in acidic aqueous solutions are strongly influenced by specific interactions with molecules in solvent environment. In aqueous solution, they found that using DFT (ωB97X-D/6-31G*) with implicit solvation (PCM(Bondi)) for geometry optimizations yields good results, except for the redox couples that have strong intramolecular hydrogen bond interactions. They evaluated a total of 19 Anthraquinones and reported a mean absolute deviation (MAD) of 0.194 V for three outliers that showed strong intramolecular hydrogen bond interactions. This value was more than five times the MAD value of the remaining 16 redox couples (0.037 V). Further, they showed that QM/MM calculations (with the TIP3P force field for explicit water molecules) alleviate the overestimation and lead to a more balanced treatment of solute–solvent interactions. Accordingly, using a QM/MM model, the correlation between theory and experiment data had a MAD of 0.033 V. In the current work, we performed a similar analysis on our calibration data set of 43 molecule pairs. We found that there are only 12 molecules, with IDs: 1, 2, 3, 4, 5, 6, 8, 9, 16, 35, 37 and 39 shown in Supplementary Table S1, without any possibility of strong intramolecular hydrogen bond interactions due to the neighboring positions found in the hydroquinone versions of the molecules. Surprisingly, we found that when using implicit aqueous solvation during geometry optimizations, the MADs for the predicted redox potentials were 0.039 V for the 12 molecules and 0.037 V for the remaining 31 molecules. We note that these MADs are very similar and the difference between the two groups is only in the third decimal digit. Therefore, we cannot confirm that the explanation provided by Kim et al. also applies to the methods used in this work. At the same time, it must be noted that Kim et al. used only Anthraquinones (3-ring molecules) for their analysis, whereas this work considers a wide variety of quinone molecules (from 1 to 3 rings), including those with the C=O groups at the 1,2 positions on the compounds. Further, Kim et al. employed the PCM (Bondi) implicit solvation model, which is different from the PBF model used in the current work. These differences, as well as the difference in the calibration data, make it hard to ascertain the exact origin of the disparities between this work and the work of Kim et al.

Comparison of low-level methods: FF, SEQM and DFTB. After establishing the effectiveness of DFT-based methods, we next consider the computationally less costly methods for the optimization of geometries and the prediction of energies. As summarized in Fig. 3 and Supplementary Fig. S4, we employ different lower level methods such as FF, SEQM and DFTB for geometry optimizations. To obtain $\Delta E_{\text{rxn}}$ of the reactions, we use the optimized structures of compounds and perform SPE calculations via the following three schemes:
I. SPE values are taken directly from low-level methods after the geometry optimizations in gas- or aqueous-phases. Aqueous-phase optimizations are performed only by using FF and SEQM (but not DFTB, as explained in Computational Workflow).

II. SPE values are taken from gas-phase DFT calculations using three different functionals (PBE, B3LYP and M08-HX), on the molecular geometries obtained via scheme (I).

III. SPE values are taken from DFT calculations with implicit solvation using three different functionals (PBE, B3LYP and M08-HX), on the molecular geometries obtained via scheme (I).

Several observations are made by comparing the $R^2$ and RMSE data across the various combinations of methods.

**Comparisons within scheme (I).** When comparing redox potential predictions from SPE data of scheme (I) to $\text{PBE}_g$ benchmark ($R^2 = 0.954$, RMSE = 0.072 V), we make the following observations (note: subscript 'g' represents gas-phase and subscript 'aq' represents aqueous-phase geometry optimizations at a given level of theory):

- In Fig. 3a, gas-phase (OPLS3eg: $R^2 = 0.596$, RMSE = 0.213 V) and aqueous-phase (OPLS3eq: $R^2 = 0.325$ V) calculated FF SPEs are significantly worse than $\text{PBE}_g$. Just as observed for DFT methods, aqueous-phase FF optimizations yield worse results than their gas-phase counterpart. Clearly, the internal energy predictions at the FF level are inaccurate.

- In Fig. 3a, gas-phase SEQM methods show significantly better performance compared to FF method, and are close to $\text{PBE}_g$ benchmark. Of note are the AM1eg ($R^2 = 0.899$, RMSE = 0.107 V) and PM7eg ($R^2 = 0.906$, RMSE = 0.103 V) methods. The aqueous-phase SEQM geometry optimizations with COSMO solvation model result in very similar prediction accuracies to their gas-phase counterparts for both AM1eq ($R^2 = 0.886$, RMSE = 0.113 V) and PM7eq ($R^2 = 0.915$, RMSE = 0.098 V).
In Fig. 3a, gas-phase DFTB methods perform as good as PBE benchmark, with parameter sets DFTB-D3 \(_g\) (R\(^2 = 0.953\), RMSE = 0.072 V) and GFN1-XTB \(_g\) (R\(^2 = 0.944\), RMSE = 0.079 V).

**Comparisons within scheme (II).** When comparing predictions from SPE data of scheme (II) to PBE \(_g\) benchmark (R\(^2 = 0.954\), RMSE = 0.072 V), we make the following observations:

- In Fig. 3b (solid bars), the performances of gas-phase DFT calculations of SPEs on gas-phase FF geometries (OPLS3\(_{eq}\); R\(^2 = 0.947\), RMSE = 0.077 V) are significantly better than their counterparts from scheme (I). The same is also observed for gas-phase DFT calculations of SPEs on aqueous-phase FF geometries (OPLS3\(_{eq}\); R\(^2 = 0.939\), RMSE = 0.083 V). However, even after performing DFT calculations of SPEs, the OPLS3\(_{eq}\) performs worse than OPLS3\(_{eq}\).

- In Fig. 3b (solid bars), gas-phase DFT calculations of SPEs on gas-phase SEQM geometries also show improved prediction accuracies with respect to their counterparts from scheme (I). The two best SEQM methods are AM1 \(_g\) (R\(^2 = 0.963\), RMSE = 0.064 V) and PM7 \(_g\) (R\(^2 = 0.954\), RMSE = 0.072 V), with performances equivalent to PBE \(_g\) benchmark. Gas-phase DFT calculations of SPE on aqueous-phase SEQM geometries resulted in worse predictions for both AM1 \(_aq\) (R\(^2 = 0.956\), RMSE = 0.070 V) and PM7 \(_aq\) (R\(^2 = 0.943\), RMSE = 0.080 V), though they are still better with respect to their counterparts from scheme (I).

- In Fig. 3b (solid bars), gas-phase DFT calculations of SPEs on gas-phase DFTB geometries also show slightly improved prediction accuracies and are slightly better than PBE \(_g\) benchmark, with parameter sets DFTB-D3 \(_g\) (R\(^2 = 0.960\), RMSE = 0.067 V) and GFN1-XTB \(_g\) (R\(^2 = 0.949\), RMSE = 0.075 V).

**Comparisons within scheme (III).** Upon including implicit solvation effects during DFT calculations of SPE, the performances of low-level methods versus the corresponding PBE \(_g\) benchmark (R\(^2 = 0.977\), RMSE = 0.051 V) can be described as follows (please note that the subscript ‘s’ represents gas-phase geometry optimization but with implicit solvation included while calculating SPE with DFT):

- In Fig. 3b (dashed bars), RMSEs are lowered by 0.02 V for FF optimized geometries, both from gas- (OPLS3\(_{eq}\); R\(^2 = 0.969\), RMSE = 0.059 V) and aqueous-phases (OPLS3\(_{eq}\); R\(^2 = 0.964\), RMSE = 0.063 V), in comparison to their counterparts from scheme (II). Surprisingly, the performances of these methods are close to PBE \(_g\) benchmark. This shows that even though the thermochemistry with FF obtained energies is not accurate (as observed in Scheme I), the quinone molecule geometries from FF are good enough for performing DFT SPE calculations.

- In Fig. 3b (dashed bars), prediction accuracies from SEQM optimized geometries are improved when compared to their counterparts in scheme (II). The two best SEQM methods are AM1 \(_s\) (R\(^2 = 0.969\), RMSE = 0.059 V) and PM7 \(_s\) (R\(^2 = 0.976\), RMSE = 0.051 V). Yet again, predictions from aqueous-phase SEQM geometries resulted in slightly worse results for both AM1 \(_aq\) (R\(^2 = 0.961\), RMSE = 0.066 V) and PM7 \(_aq\) (R\(^2 = 0.962\), RMSE = 0.065 V). Interestingly, the performances of these SEQM methods are close to PBE \(_g\) benchmark.

- In Fig. 3b (dashed bars), prediction accuracies from DFTB optimized geometries improve when compared to their counterparts in scheme (II). Strikingly, both sets of DFTB parameters, DFTB-D3 \(_s\) (R\(^2 = 0.978\), RMSE = 0.049 V) and GFN1-XTB \(_s\) (R\(^2 = 0.977\), RMSE = 0.051 V), perform better than PBE \(_g\) benchmark.

**Discussion**

All the variations in computational methods that are used for geometry optimizations and SPE calculations, also with and without implicit solvation effects, are found to influence the prediction accuracies to varying degrees. First, for all methods, gas-phase DFT calculations of SPEs lead to significant improvements in prediction accuracies. According to these results, computationally demanding DFT geometry optimizations are hardly necessary for a first-order screening of large numbers of candidate molecules. Instead, either of SEQM or DFTB methods may be employed for the task of gas-phase geometry optimizations. Secondly, SPE calculations employing the PBE functional are generally better performing than SPEs obtained from computationally more costly B3LYP and M08-HX functionals. Thirdly, for all low-level methods, the inclusion of implicit solvation during DFT calculations of SPEs leads to improved prediction accuracies. Finally, the results confirm that the effects of geometry optimizations in aqueous-phase are minimal and they often result in slightly worse prediction accuracies.

The calibration equation for the prediction of redox potentials versus SHE, using the DFTB-D3 \(_g\) (\(\equiv T\)) calculated reaction energies is:

\[
E^0 = -0.447[\Delta E_{\text{DFTB}}] - 0.823
\]  

(2)

The performance metrics of all the low-level methods, as well as their combinations with high-level methods, are given in Supplementary Tables S3–S9. It is worth pointing out that a recent study by Tabor et al.\(^{27}\) also considered DFT and SEQM based methods for performing a similar calibration exercise to determine an optimum method for predicting redox potentials of 28 quinone molecules using \(\Delta E_{\text{DFTB}}\) as a descriptor. In their study, at the DFT level, they considered the B3LYP functional with 6-31+G(d,p) basis set, which is very similar to the basis set used in this study. At the SEQM level, they considered the PM7 method. The MAD with gas-phase DFT was 0.056 V (vs. 0.058 V using PBE in this work) and with implicit aqueous solvation using the PCM model was 0.041 V (vs. 0.038 V using PBE with PBF solvation model in this work). When using the PM7 method, their predicted MAD value was 0.052 V in the gas phase (vs. 0.065 V in this work), and 0.067 V with the COSMO solvation model (vs. 0.067 V in this work). The difference between the results from the two studies are quite small.
and vary mostly in the third decimal place, which is not surprising given the common pool of molecules. For predicting redox potentials, we recommend using Eq. 1 not only because it has slightly better prediction accuracy but also because it covers a larger group of experimented molecules with a wider range of redox potentials.

**Accuracy of predictions versus cost of calculations.** In addition to determining accurate methods for predicting the redox potentials of electroactive quinones, a major aim of the current study is to decide on the methods that are most suited for both standalone and HTCS studies for which a balance between the speed of computations and the accuracy of results is desirable. This is especially important when the robust DFT calculations become impractical for studies on a vast chemical space (10^3 ~ 10^6 compounds). Thus, an estimation of trade-offs between the computational accuracy and its cost is useful for efficient screening studies. For a comparison of the computation time, we selected a representative method for each level of theory, namely, OPLS3e (FF), PM7 (SEQM), DFTB-D3 (DFTB) and PBE (DFT). Next, noting that the geometry optimizations are usually the most computationally demanding processes, we optimized the geometries of all the molecules in gas-phase using these representative methods. We added the FF geometry optimization time to all other methods’ calculation times, since we use it as the base method for performing all other geometry optimizations (as explained in Computational Workflow). An averaged computation time, as obtained from five different runs that actually showed no significant variation, is used to describe the relation between RMSEs versus the computational cost. As shown in Fig. 4a, DFTB-D3 is almost as accurate as DFT in predicting the redox potentials, while it requires substantially less (~ 10^3 times) computing time. Also noting that DFTB-D3 has previously been applied for calculations of large systems at relatively lower computational costs and with similar accuracies to that of the higher level (i.e. DFT-GGA) methods, we suggest that for the small quinone redox compounds the DFTB-D3 method provides a good compromise between prediction accuracy and computational cost. Accordingly, for HTCS studies that are aimed to work on extremely large chemical spaces of molecules, we suggest DFTB-D3 computations on OPLS3e optimized geometries, which have an RMSE of 0.072 V that is close to the RMSE of 0.051 V of the DFT (PBE) calculations, as a way to accelerate the virtual screening of compounds.

**Effects of geometry optimizations at various levels of theory and implicit solvation.** For the set of methods considered in the previous section, we also quantify the effect of gas-phase DFT calculation of SPEs by using the geometries that have been obtained via low-level methods. As shown in Fig. 4b (solid bars), the improvement in prediction accuracy, Δe RMSE, is most significant for geometries optimized by OPLS3e (Δe RMSE = 0.136 V), followed by PM7 (Δe RMSE = 0.031 V), and then by DFTB-D3 (Δe RMSE = 0.005 V). These results show that SEQM, and more pronouncedly, DFTB methods do not only predict the reaction energies accurately, but they also predict the geometries of the compounds as comparable to that of DFT. For the same set of methods, we also investigate relationships between the differences in molecular geometries and the calculated values of Δe RMSE. By performing structure superposition analysis we compared the optimized geometries from different methods to the reference, gas-phase PBE optimized geometries. The average root-mean-square deviation (RMSD) of all the 86 reactant and product molecules under various atomic constraints are shown in Table 1. First, under all constraints OPLS3e has the largest average RMSD with respect to PBE, which is expected. Secondly, when considering all atoms or only heavy atoms, PM7 and DFTB-D3 are very similar in geometrical difference with regards to PBE. Thirdly, it is surprising to note that DFTB-D3, while the most accurate of the low-level methods, does not necessarily provide the geometry closest to PBE (i.e. lowest average RMSD) when
all atoms or non-hydrogen atoms are considered. However, when considering only the carbon atoms in the ring structure of molecules, DFTB-D3 produces structures that are closest to PBE. Given that the cyclic carbon atoms are a large fraction of the total number of atoms, it is possible that being able to represent the geometry of the rings accurately is what gives DFTB-D3 (Δe RMSE = 0.005 V) an advantage over PM7 (Δe RMSE = 0.031 V) and OPLS3e (Δe RMSE = 0.136 V) for the redox potential predictions.

Next, we quantify the improvements in prediction capability of the methods, with respect to gas-phase DFT calculation of SPE, due to the inclusion of implicit solvation effects, Δe RMSE (dashed bars). As shown in Fig. 4b, an improvement in prediction accuracy is evident for all levels of theory and Δe RMSE are similar at each level of theory. The decrease in RMSEs are 0.018, 0.021, 0.018, and 0.021 V for OPLS3e, PM7, DFTB-D3, and PBE, respectively. These results show that the amount of improvement due to the inclusion of implicit solvation is independent of the source of geometry. These findings are useful, for instance, when building machine learning (ML) models for the prediction of solvation energies directly from simple cheminformatics-based descriptors without a real need for an explicit knowledge of compound geometries.

In addition to the above findings, the results presented here are expected to be useful for generating accurate and large quantity chemical data on compounds, which can later be utilized by data-driven machine learning models for expanding the boundaries of search space during candidate compound explorations. Firstly, as has been argued in recent studies27,29, using the computationally costly quantum chemical calculations for millions of molecules, which is required for building powerful ML models, is still a major bottleneck. One of the key findings of this work is that the DFTB method is nearly as accurate as DFT when it comes to the prediction of quinone redox potentials. Thus, the data scarcity bottleneck can be addressed directly by using DFTB to generate large quantities of reliable reference data. Secondly, it has been demonstrated quantitatively in this work that the improvement due to the use of quantum chemistry methods, when compared to SEQM and DFTB methods, corresponds to an energy contribution that constitutes only a minor fraction of the total energy. An increasingly popular strategy that utilizes this fact for the property predictions at quantum chemical accuracy is Δ-ML42. In this strategy, by training on quantum chemical reference data, a ML model is trained to predict the correction terms, Δ, to values that have been computed using the inexpensive methods. This way, instead of performing quantum chemical calculations directly on all candidate compounds, the energy values are interpolated from a selected inexpensive baseline theory to quantum chemical accuracy. In our study we recommend a choice of methods, for both reference (e.g. M08-HX) and base line (e.g. DFTB-D3) data, to perform such Δ-ML tasks. Finally, we would like to emphasize again, that this work provides a general methodology for determining the optimum combination of methods with an example case of the quinone family of compounds. This knowledge is directly relevant for developing advanced decision-making frameworks that can automatically decide on the most optimum combination of computational methods for a target class of chemical compounds.

Methods

Thermodynamic principle. The reaction energy difference between the reactant and product is used as an approximation for the Gibbs free energy of proton-coupled electron transfer redox reactions. During a redox reaction in the aqueous-phase:

$$Q + 2H^+ + 2e^- \rightarrow QH_2$$

The hydroquinone, QH₂, compounds can be generated from the quinone, Q, compounds via a two-electron two-proton redox reaction22. A quantitative measure of the favorability of a given reaction is the change in standard Gibbs free energy, ΔG°(rxn). According to the Nernst equation, the equilibrium potential of a redox reaction, E°, is related to the change in the standard Gibbs free energy per coulomb of charge transferred during the electrochemical reaction as:

$$E^o = -\frac{\Delta G^o_{\text{rxn}}}{nF}$$

where n = 2 is the number of electrons and F is the Faraday constant. Typically, E° is measured relative to the Standard Hydrogen Electrode (SHE) and ΔG°(rxn) is computed at standard conditions. To calculate ΔG°(rxn) we use the following equation:

$$\Delta G^o_{\text{rxn}} = \Delta G^o(QH_2) - [G^o(Q) + G^o(H_2)]$$

in which ΔG°(rxn) is expressed simply as the difference in the standard free energies of the reactants and products. Thermodynamically, ΔG°(rxn) can be described as a sum of contributions arising from the change in internal energy (ΔUrxn), pressure–volume (pΔVrxn) and entropic (TΔSrxn) contributions due to reaction as:

| Method          | All atoms | Non-hydrogen | Only ring carbon |
|-----------------|-----------|--------------|-----------------|
| OPLS3e (FF)     | 0.148     | 0.120        | 0.052           |
| PM7 (SEQM)      | 0.132     | 0.098        | 0.050           |
| DFTB-D3 (DFTB)  | 0.135     | 0.102        | 0.043           |

Table 1. The difference in the optimized molecular geometry using various calculation methods with reference to PBE (DFT) geometry under various atomic constraints. The average RMSD values for the 86 molecules are shown in units of Å. RMSD values have been calculated for all atoms, all non-hydrogen atoms, and carbon atoms of molecule rings.
\[ \Delta G_{\text{rxn}}^o = \Delta U_{\text{rxn}} + p \Delta V_{\text{rxn}} - T \Delta S_{\text{rxn}} \]  

(6)

The change in internal energy can be further decomposed as \( \Delta U_{\text{rxn}} = \Delta E_{\text{rxn}} + \Delta ZPE \), where \( \Delta E_{\text{rxn}} \) is the reaction energy and \( \Delta ZPE \) is the change in zero-point energy. In the present work, the ZPE contributions to the internal energy, changes in pressure–volume, and entropic contributions are neglected, thus effectively using the approximation:

\[ \Delta G_{\text{rxn}}^o \cong \Delta E_{\text{rxn}} = E^{\text{q}}(\text{QH}_2) - \left[ E^{\text{q}}(\text{Q}) + E^{\text{q}}(\text{H}_2) \right] \]  

(7)

where \( E^{\text{q}} \) represents the theoretically calculated internal energy of species in aqueous-phase. All the other contributions are ignored because they require extra calculation steps, computational resources, and thus, are not suitable from the perspective of HTCS. Nevertheless, the effects of ignoring these other contributions on the accuracy of predictions are discussed in the main text. In this work, the descriptor of choice, \( \Delta E_{\text{rxn}} \), is calculated with the inclusion of aqueous solvation effects, which requires additional computation time. To quantify the effect of solvation on prediction accuracy, another approximation is considered by ignoring solvation such that Eq. (7) can be rewritten using internal energies calculated in the gas-phase as:

\[ \Delta G_{\text{rxn}}^o \cong \Delta E_{\text{rxn}} = E^{\text{g}}(\text{QH}_2) - \left[ E^{\text{g}}(\text{Q}) + E^{\text{g}}(\text{H}_2) \right] \]  

(8)

Under these set of approximations, the calculated change in internal energies \( \Delta E_{\text{rxn}} \) from Eqs. (7) and (8) are linked to the measured redox potentials using Eq. (4). We used various theoretical methods, as explained below in computational details, to calculate \( \Delta E_{\text{rxn}} \), and discuss their performance for the prediction of experimentally measured redox potentials.

Computational details. We developed a systematic computational approach involving one FF, nine SEQM, two DFTB, and eleven DFT methods, as well as their combination with implicit solvation environments, to predict the experimentally measured redox potentials of quinone-based electroactive. The MacroModel program is used for FF configurational searches and geometry optimizations, while the Jaguar program is used for DFT calculations, all as implemented in the Schrödinger Materials Science Suite (version 2019-2). MOPAC and DFTB calculations are performed using the ADF program. Molecular structures are optimized both in gas- and aqueous-phase using the OPLS3e FF, which provides a broad coverage of small compounds. The gas-phase FF optimized geometries are used as inputs to perform gas- and aqueous-phase geometry optimizations using nine different SEQM methods, including AM1, MNDO, MNDO-D, PM3, PM6, PM6-D3, PM6-D3H4X, PM7, and RM1. The aqueous-phase geometry optimizations at the SEQM level are performed using the COSMO-RS solvation model. The choice of this solvation method is constrained by the present availability in the ADF program. The gas-phase FF optimized geometries are also used as inputs for DFTB optimizations using the DFTB-D3 and GFN1-xTB methods. The DFTB-D3 computations are performed with a self-consistent charge cycle using the QuasiNANO-2015 parameter set, while the parameters for GFN1-xTB are taken from the work of Grimme et al. The aqueous-phase geometry optimizations of molecules are not performed with the DFTB method, since currently there is no available routine for this task in the ADF program. Finally, FF minimized geometries are used as inputs to perform geometry optimizations in gas-phase FF calculations using different flavors of exchange–correlation functionals, including the local density approximation (LDA), generalized gradient approximation (GGA), and hybrid meta-GGA functionals, all of which vary drastically in their accounting of the exchange–correlation energy. For the geometries that have been obtained from FF, SEQM and DFTB optimizations, the DFT level SPEs are computed in gas-phase, and subsequently in aqueous-phase using only the PBE, B3LYP, and M08-HX functionals, as they are well accepted in the community but also span a wide range of ways for accounting for the exchange–correlation effects. A total of 11 exchange–correlation functionals, also including some of the D3 dispersion corrected variants, are used for high-level OPT and SPE calculations. These functionals include LDA, PBE, PBE-D3, BLYP, BLYP-D3, B3LYP, B3LYP-D3, BHLYP, HSE06, and M08-HX. Due to computational costs, the DFT aqueous-phase geometry optimizations are performed only with the following functionals: PBE, B3LYP, and M08-HX.

As DFT options in Jaguar, we chose “medium” grid density for OPT and “fine” grid density for SPE calculations. Energy and RMS density matrix change convergence criteria are set to the default values of 5.0 × 10⁻⁵ and 5.0 × 10⁻⁶ Hartree, respectively. The default direct inversion in the iterative subspace is employed as the convergence scheme. For OPT, Jaguar’s mixed pseudospectral grids with default cutoffs are employed. For SPE calculations, we used pseudospectral grids with accurate cutoffs. To treat solvated molecules in water, we used the standard PBF solver with water as the solvent. The calculations are performed with LACVP⁺⁺ basis set with polarization and diffuse functions. The LACVP basis set is chosen here because it includes an effective core potential (ECP), which represents the effect of the core electrons in a parametrized form. The use of ECPs speeds up calculations on compounds that contain heavy elements. For the elements from H to Ar, LACVP and the widely employed 6-31G are essentially indistinguishable when evaluating the ground-state properties. The quinone molecules considered in this work contain the elements of C, H, O, N, S, F, and Cl, and thus, the use of LACVP⁺⁺ basis set in this work is consistent with the use of 6-31G⁺⁺ basis set.

Calibration data and performance metrics. We collected redox potential data from 43 quinone redox couples in acidic aqueous solution. In consideration of prediction accuracy and universality for the calibration models, the selection of available experimental data has been expanded within various quinone molecules, rather than using monotonous structural patterns. This is because compounds decorated with chemical functional groups usually show different redox potentials as well as charge/discharge capacities when compared
to their undecorated counterparts. The experimental molecules cover both quinone cores and their functionalized derivatives with multiple substituted groups including –SO$_3$H, –COOH, –OH, –CH$_3$, –F and –Cl (see Supplementary Table S1) and the redox data spans a broad range experimental redox potentials between −0.084 and 1.21 V. The redox couples are chosen consistently from measurements that were performed under similar experimental conditions, such as $T = 298.15$ K, pH = 0, and highly conducting salts. The correlations between experiments and calculations are expressed in terms of two commonly used coefficients, namely the coefficient of determination ($R^2$) and root-mean-square error (RMSE). $R^2$ and RMSE are calculated using the definitions from the Originlab. All benchmark simulations of calculation time have been performed on a single core of Intel Core i9-9960X 3.10 GHz CPU with Ubuntu 18.04 Bionic Beaver as the operating system.

Data availability
The generated computational data of compounds is provided in Supplementary Table S1.

Received: 11 September 2020; Accepted: 2 December 2020
Published online: 17 December 2020

References
1. Ding, Y., Zhang, C., Zhang, L., Zhou, Y. & Yu, G. Molecular engineering of organic electroactive materials for redox flow batteries. Chem. Soc. Rev. 47, 69–103 (2018).
2. Singh, V., Kim, S., Kang, J. & Byon, H. R. Aqueous organic redox flow batteries. Nano Res. 12, 1–14 (2019).
3. Nakai, Y., Li, Y. & Aziz, M. J. Electrolyte lifetime in aqueous organic redox flow batteries: A critical review. Chem. Rev. 120, 6467–6489 (2020).
4. Noack, J., Roznyatovskaya, N., Herr, T. & Fischer, P. The chemistry of redox-flow batteries. Angew. Chem. Int. Ed. 54, 9776–9809 (2015).
5. Webers, A. Z. et al. Redox flow batteries: A review. J. Appl. Electrochem. 41, 1137–1164 (2011).
6. Kears, G., Shah, A. A. & Walsh, F. C. Development of the all-vanadium redox flow battery for energy storage: A review of technological, financial and policy aspects. Int. J. Energy Res. 36, 1105–1120 (2012).
7. Winsberg, J., Hagemann, T., Janoschka, T., Hager, M. D. & Schubert, U. S. Redox-flow batteries: From metals to organic redox-active materials. Angew. Chem. Int. Ed. 55, 2–28 (2016).
8. Debrulier, C., Hu, B., Mose, J., Luo, J. & Liu, T. L. A sulfonate-functionalized viologen enabling neutral cation exchange, aqueous organic redox flow batteries toward renewable energy storage. ACS Energy Lett. 3, 663–668 (2018).
9. Liu, W. et al. A highly stable neutral viologen/bromine aqueous flow battery with high energy and power density. Chem. Commun. 55, 4801–4804 (2019).
10. Lin, K. et al. A redox-flow battery with an alloxazine-based organic electrolyte. Nat. Energy 1, 16102 (2016).
11. Lee, W., Kwon, B. W. & Kwon, Y. Effect of carboxylic acid-doped carbon nanotube catalyst on the performance of aqueous organic redox flow battery using the modified alloxazine and ferrocyanide redox couple. ACS Appl. Mater. Interfaces 10, 36882–36891 (2018).
12. Hollas, A. et al. A biomimetic high-capacity phenazine-based anolyte for aqueous organic redox flow batteries. Nat. Energy 3, 508–514 (2018).
13. Winsberg, J. et al. TEMPO/phenazine combi-molecule: A redox-active material for symmetric aqueous redox-flow batteries. ACS Energy Lett. 1, 976–980 (2016).
14. Orita, A., Verde, M. G., Sakai, M. & Meng, Y. S. The impact of pH on side reactions for aqueous redox flow batteries based on nitroxyl radical compounds. J. Power Sources 321, 126–134 (2016).
15. Li, L. et al. High-capacity aqueous zinc batteries using sustainable quinone electrodes. Sci. Adv. 4, 1761 (2018).
16. Gerhardt, M. K. et al. Anthraquinone derivatives in aqueous flow batteries. Adv. Energy Mater. 7, 1601488 (2017).
17. Wang, C. et al. High-performance alkaline organic redox flow batteries based on 2-hydroxy-3-carboxy-1,4-naphthoquinone. ACS Energy Lett. 3, 2404–2409 (2018).
18. Lee, B. et al. Exploiting biological systems: Toward eco-friendly and high-efficiency rechargeable batteries. Joule 2, 61–75 (2018).
19. Wedge, K., Dražević, E., Konya, D. & Bentien, A. Organic redox species in aqueous flow batteries: Redox potentials, chemical stability and solubility. Sci. Rep. 6, 1–13 (2016).
20. Goulet, M.-A. et al. Alkaline benzoquinone aqueous flow battery for large-scale storage of electrical energy. Adv. Energy Mater. 8, 1702056 (2017).
21. Ji, Y. et al. A Phosphonate-functionalized quinone redox flow battery at near-neutral pH with record capacity retention rate. Adv. Energy Mater. 1900039, 1–7 (2019).
22. Huskinson, B. et al. A metal-free organic-inorganic aqueous flow battery. Nature 505, 195–198 (2014).
23. Yang, B., Narayanan, S. R., Hoober-Burkhardt, L., Wang, F. & Surya Prakash, G. K. An inexpensive aqueous flow battery for large-scale electrical energy storage based on water-soluble organic redox couples. J. Electrochem. Soc. 161, 1371–1380 (2014).
24. Sorkun, M. C., Khetan, A. & Er, S. AqSolDB, a curated reference set of aqueous solubility and 2D descriptors for a diverse set of compounds. Sci. data 6, 143 (2019).
25. Cheng, L. et al. Accelerating electrolyte discovery for energy storage with high-throughput screening. J. Phys. Chem. Lett. 6, 283–291 (2015).
26. Er, S., Suh, C., Marshak, M. P. & Aspuru-Guzik, A. Computational design of molecules for an all-quinone redox flow battery. Chem. Sci. 6, 885–893 (2015).
27. Tabor, D. P. et al. Mapping the frontiers of quinone stability in aqueous media: Implications for organic aqueous redox flow batteries. J. Mater. Chem. A 7, 12833–12841 (2019).
28. Huyhn, M. T., Anson, C. W., Cavell, A. C., Stahl, S. S. & Hammes-Schiffer, S. Quinone 1 e- and 2 e-/2 H+ reduction potentials: Identification and analysis of deviations from systematic scaling relationships. J. Am. Chem. Soc. 138, 15903–15910 (2016).
29. Jinich, A., Sanchez-Lengeling, B., Ren, H., Harman, R. & Aspuru-Guzik, A. A Mixed quantum chemistry/machine learning approach for the fast and accurate prediction of biochemical redox potentials and its large-scale application to 315 000 redox reactions. ACS Cent. Sci. 5, 1199–1210 (2019).
30. Inset, J.-P., Ramesh, S., Duan, C. & Kulik, H. J. Accurate multiojective design in a space of millions of transition metal complexes with neural-network-driven efficient global optimization. ACS Cent. Sci. 6, 513–524 (2020).
31. Bachmann, J. et al. Lead candidates for high-performance organic photovoltaics from high-throughput quantum chemistry-the Harvard clean energy project. Energy Environ. Sci. 7, 698–704 (2014).
32. Kim, H., Goodson, T. & Zimmerman, P. M. Achieving accurate reduction potential predictions for anthraquinones in water and aprotic solvents: Effects of inter- and intramolecular H-bonding and ion pairing. J. Phys. Chem. C 120, 22235–22247 (2016).
33. Oliveira, A. E., Philippsen, P. & Heine, T. DFTB parameters for the periodic table, part 2: Energies and energy gradients from hydrogen to calcium. *J. Chem. Theory Comput.* 11, 5209–5218 (2015).
34. Christensen, A. S., Kubat, T., Cui, Q. & Elstner, M. Semiempirical quantum mechanical methods for noncovalent interactions for chemical and biochemical applications. *Chem. Rev.* 116, 5301–5337 (2016).
35. Weininger, D. SMILES, a chemical language and information system. 1. Introduction to methodology and encoding rules. *J. Chem. Inf. Comput. Sci.* 1, 31–36 (1988).
36. Roos, K. et al. OPLSae: Extending force field coverage for drug-like small molecules. *J. Chem. Theory Comput.* 15, 1863–1874 (2019).
37. Tannor, D. J. et al. Accurate first principles calculation of molecular charge distributions and solvation energies from ab initio quantum mechanics and continuum dielectric theory. *J. Am. Chem. Soc.* 116, 11875–11882 (1994).
38. Bochevarov, A. D. et al. Jaguar: A high-performance quantum chemistry software program with strengths in life and materials sciences. *Int. J. Quantum Chem.* 113, 2110–2142 (2013).
39. Mardirossian, N. & Head-Gordon, M. How accurate are the minnesota density functionals for noncovalent interactions, isomerization energies, thermochemistry, and barrier heights involving molecules composed of main-group elements? *J. Chem. Theory Comput.* 12, 4303–4325 (2016).
40. Medvedev, M. G., Bushmarinov, I. S., Sun, J., Perdew, J. P. & Lyssenko, K. A. Density functional theory is straying from the path toward the exact functional. *Science* 355, 49–52 (2017).
41. Zheng, G. et al. Parameter calibration of transition-metal elements for the spin-polarized self-consistent-charge density-functional tight-binding (DFTB) Method: Sc, Ti, Fe, Co, and Ni. *J. Chem. Theory Comput.* 3, 1349–1367 (2007).
42. Ramakrishnan, R., Dral, P. O., Rupp, M. & Von Lilienfeld, O. A. Big data meets quantum chemistry approximations: The Δ-machine learning approach. *J. Chem. Theory Comput.* 11, 2087–2096 (2015).
43. Te Velde, G. et al. Chemistry with ADF. *J. Comput. Chem.* 22, 931–967 (2001).
44. Dewar, M. J. S., Zoebisch, E. G., Healy, E. F. & Stewart, J. J. P. AM1: A new general purpose quantum mechanical molecular model. *J. Am. Chem. Soc.* 107, 3902–3909 (1985).
45. Dewar, M. J. S. & Thiel, W. Ground states of molecules. 38. The MND0 method. Approximations and parameters. *J. Am. Chem. Soc.* 99, 4899–4907 (1977).
46. Thiel, W. & Voityuk, A. A. Extension of MNDO to d orbitals: Parameters and results for the second-row elements and for the zinc group. *J. Phys. Chem.* 100, 616–626 (1996).
47. Stewart, J. J. P. Optimization of parameters for semiempirical methods II applications. *J. Comput. Chem.* 10, 221–264 (1989).
48. Stewart, J. J. P. Optimization of parameters for semiempirical methods V: Modification of NDDO approximations and application to 70 elements. *J. Mol. Model.* 13, 1173–1213 (2007).
49. Jeanvoine, Y. & Spezia, R. The formation of urea in space. II. MP2 versus PM6 dynamics in determining bimolecular reaction products. *Theor. Chem. Acc.* 138, 1–13 (2019).
50. Stewart, J. J. P. Optimization of parameters for semiempirical methods VI: More re-optimization of parameters. *J. Mol. Model.* 19, 1–32 (2013).
51. Rocha, G. B., Freire, R. O., Simas, A. M. & Stewart, J. J. P. RM1: A reparametrization of AM1 for H, C, N, O, P, S, F, Cl, Br, and I. *J. Comput. Chem.* 27, 1101–1111 (2006).
52. Klamt, A. & Eckert, F. COSMO-RS a novel and efficient method for the a priori prediction of thermophysical data of liquids. *Fluid Phase Equilib.* 172, 43–72 (2000).
53. Chu, Y., Zhang, X., Hillestad, M. & He, X. Computational prediction of cellulose solubilites in ionic liquids based on COSMO-RS. *Fluid Phase Equilib.* 475, 25–36 (2018).
54. Brandenburg, J. G. & Grimme, S. Accurate modeling of organic molecular crystals by dispersion-corrected density functional tight binding (DFTB). *J. Phys. Chem. Lett.* 5, 1785–1789 (2014).
55. Grimme, S., Bannwarth, C. & Shushkov, P. A robust and accurate tight-binding quantum chemical method for structures, vibrational frequencies, and noncovalent interactions of large molecular systems parametrized for all spd-block elements (Z = 1–86). *J. Chem. Theory Comput.* 13, 1809–2007 (2017).
56. Koopman, J. & Grimme, S. Calculation of electron ionization mass spectra with semiempirical GFN-xTB methods. *ACS Omega* 4, 15120–15133 (2019).
57. Vosko, S. H., Wilk, L. & Nusair, M. Accurate spin-dependent electron liquid correlation energies for local spin density calculations: A critical analysis. *Can. J. Phys.* 58, 1200–1211 (1980).
58. Mardirossian, N. & Head-Gordon, M. Thirty years of density functional theory in computational chemistry: An overview and extensive assessment of 200 density functionals. *Mol. Phys.* 115, 2315–2372 (2017).
59. Swart, M. A new family of hybrid density functionals. *Chem. Phys. Lett.* 580, 166–171 (2013).
60. Hujo, W. & Grimme, S. Comparison of the performance of dispersion-corrected density functional theory for weak hydrogen bonds. *Phys. Chem. Chem. Phys.* 13, 13942–13950 (2011).
61. Perdew, J. P., Burke, K. & Ernzerhof, M. Generalized gradient approximation made simple. *Phys. Rev. Lett.* 78, 3865–3868 (1996).
62. Stephens, P. J., Devlin, F. J., Chabalowski, C. F. & Frisch, M. J. Ab initio calculation of vibrational absorption and circular dichroism spectra using density functional force fields. *J. Phys. Chem.* 98, 11623–11627 (1994).
63. Sánchez-Castellanos, M. et al. Theoretical exploration of 2,2'-bipyridines as electro-active compounds in flow batteries. *Phys. Chem. Chem. Phys.* 21, 15823–15832 (2019).
64. Adamo, C. & Barone, V. Toward reliable density functional methods without adjustable parameters: The PBE0 model. *Chem. Phys.* 110, 6158–6170 (2001).
65. Krukau, A. V., Vydrov, O. A., Izmaylov, A. F. & Scuseria, G. E. Influence of the exchange screening parameter on the performance of screened hybrid functionals. *J. Chem. Phys.* 125, 224106 (2006).
66. Hay, P. J. & Wadt, W. R. Ab initio effective core potentials for molecular calculations. Potentials for K to Au including the outermost core orbitale. *J. Chem. Phys.* 82, 299–310 (1985).
67. Jaguar, version 10.5. Schrodinger, Inc., New York, (2019).
68. Tobias Johnsson Wass, J. R., Ahlberg, E., Panas, I. & Schiffrin, D. J. Quantum chemical modeling of the reduction of quinones. *J. Phys. Chem. A* 110, 2005–2020 (2006).

**Acknowledgements**

This research received funding from the Netherlands Organisation for Scientific Research (NWO), through the COLORFLOW project partnership of DIFFER and Green Energy Storage, in the framework of the Materials for Sustainability programme and from the Ministry of Economic Affairs in the framework of the “PS₃-Toezlagregeling” Grant no 739.017.013. SE acknowledges funding from the initiative “Computational Sciences for Energy Research” of Shell and NWO grant no 15CSIT105. This work was sponsored by NWO Exact and Natural Sciences for the use of supercomputer facilities.
Author contributions
S.E. conceived and directed this work. Q.Z. performed the calculations and acquired the computational data. Q.Z. and A.K. interpreted the computational data. The manuscript was written with significant contributions from all authors.

Competing interests
The authors declare no competing interests.

Additional information
Supplementary Information The online version contains supplementary material available at https://doi.org/10.1038/s41598-020-79153-w.

Correspondence and requests for materials should be addressed to S.E.

Reprints and permissions information is available at www.nature.com/reprints.

Publisher’s note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article’s Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article’s Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http://creativecommons.org/licenses/by/4.0/.

© The Author(s) 2020