1. Introduction

The majority of close-in planets found by NASA’s Kepler satellite throughout the past decade are smaller than Neptune, but larger than Earth (Batalha et al. 2013; Howard 2013; Mullally et al. 2015). The Kepler and K2 missions have shown that, of the planets to which Kepler is most sensitive ($P < 100$ days, $R_p > 1.0R_⊕$), these smaller planets are by far the most common in the galaxy (Fressin et al. 2013; Fulton et al. 2017), though there is no analog in the solar system from which this could have been predicted.

A gap in the population of planets at radii larger than 4.0 $R_⊕$ (i.e., larger than Neptune) is satisfactorily explained by runaway gas accretion (Pollack et al. 1996; Ida & Lin 2004; Mordasini et al. 2009). Larger planets are massive enough to accrete H and He from the protoplanetary disk, becoming puffy and increasing in radius. However, refined studies of the distribution of planets within the 1–4 $R_⊕$ range have revealed a significant drop in the population, or “Fulton gap” (shown in Figure 1) between 1.5 and 2.0 $R_⊕$ (Owen & Wu 2013; Fulton et al. 2017; Fulton & Petigura 2018), which is not yet well understood.

Photoevaporation presents a possible explanation for the gap, and is a particularly important factor for the close-in planets preferentially detected by Kepler. Planets with radii between 1.5 and 2.0 $R_⊕$ could represent a relatively rare group of planets retaining thin atmospheres, while super-Earths are photoevaporated rocky bodies and the sub-Neptunes are massive enough to retain thick atmospheres (Lopez & Rice 2016). Jin & Mordasini (2018) find support for this theory using planetary formation and evolution models. They observe that planets of increasing radius are more volatile-rich, with an anti-correlation between density and orbital distance. Furthermore, Fulton & Petigura (2018) find observational evidence for the photoevaporation theory in their discovery that populations of sub-Neptunes shift to higher levels of incident flux for higher-mass stars. Since stellar activity driven by rotation and convection is generally stronger and longer-lived in lower-mass stars, the atmospheres of planets orbiting smaller stars experience prolonged exposure to high-energy X-ray and UV
2. Observations and Analysis

Identified as a planet candidate from C17 of K2 (see Crossfield et al. 2018), Wolf 503 was recognized as an excellent host for follow-up study because it is both bright ($K_p = 9.9$) and nearby (45 pc). Here, we present the treatment of the photometry used to detect Wolf 503b, as well as our planet validation techniques, and derive both planetary and stellar parameters.

2.1. Photometry Extraction and Transit Detection

The photometric extraction and transit detection methods used to identify Wolf 503b are the same as those applied by our team to all light curves in C17 and are described in our corresponding C17 summary paper, Crossfield et al. (2018). As K2 operates using only two of Kepler’s four initial reaction wheels, the telescope drifts along its roll axis by a few pixels every several days, and thruster firings are used to maintain the telescope’s pointing. The change in flux resulting from this drift is removed by fitting the flux as a function of position along the drift path, which is highly similar between thruster firings. However, the data acquired during these thruster burns are not reliable and are masked out, as in the first transit of the light curve for Wolf 503, shown in Figure 2.

With the extracted light curve, we detected a candidate at $P = 6.0$ days with $S/N = 38$ having 11 transits throughout the time of observation. The candidate was marked as a particularly intriguing KOI for its favorable host star following the manual vetting procedure of the C17 candidates.

2.2. Previous Work on Wolf 503

Wolf 503 (BD-05 3763, MCC 147, LHS 2799, G 64–24, HIP 67285, TYC 4973-1501-1, 2MASS J13472346-0608121) has been a sparsely studied nearby cool star since its discovery a century ago as a high proper motion star by Wolf (1919). The star subsequently appeared in several high proper motion catalogs over the past century, as G 64–24 in Giclas et al. (1961), with Wilhelm Luyten designating the star no fewer than six times in his proper motion catalogs.35

The star was classified in numerous spectral survey as a K5V by Upgren et al. (1972), identified as UPG 336, and Bidelman (1985) published Kuiper’s posthumous classification for the star as K4 from his 1937 to 1944 survey. Pickles & Degapone (2010) found that the best-fit template for the $B_T V_T J H K_S$ photometry was that for a K4V star.

2.2.1. Distance, Kinematics, and Stellar Population

Recently, Gaia DR2 provided an ultra-precise trigonometric parallax ($\varpi = 22.430 \pm 0.048$ mas; corresponding to $d = 44.583 \pm 0.096$ pc), as well as precise proper motion and radial velocity measurements, which are listed in Table 1. Gaia itself measured a radial velocity of $-46.64 \pm 0.50 \text{ km s}^{-1}$ (2 observations), and independently, Sperauskas et al. (2016) reported a radial velocity of $-47.4 \pm 0.7 \text{ km s}^{-1}$ based on 2 CORAVEL measurements over 98 days. Combining the Gaia DR2 position, proper motion, and parallax, and the mean Gaia DR2 ground-based radial velocity (from HARPS), we estimate

35 Entry #402 in Luyten (1923) (stars with motions exceeding 0.5 5 yr$^{-1}$), as LPM 492 in Luyten (1941), LFT 1037 in Luyten (1955), LHS 2799 in Luyten (1979), and as NLTT 35228 and LTT 5351 in Luyten (1980).
Figure 2. Extracted light curve for Wolf 503 (EPIC 212779563). Transit times according to our fit are indicated with a red line. The first observed transit is not easily visible in this plot because the transit coincided with a thruster burn during which two data points were flagged and removed (see Figure 6).

### Table 1: Stellar Parameters

| Parameter | Value | Source |
|-----------|-------|--------|
| Identifying Information | | |
| α R.A. (hh:mm:ss) J2000 | 13:47:23.4439 | (Millionard 1987) |
| δ Decl. (dd:mm:ss) J2000 | −06:08:12.731 | (Millionard 1987) |
| EPIC ID | 212779563 | Gaia DR1 |
| Photometric Properties | | |
| B (mag) | 11.30 ± 0.01 | (Millionard 1987) |
| V (mag) | 10.28 ± 0.01 | (Millionard 1987) |
| G (mag) | 9.808 ± 0.001 | Gaia DR1 |
| J (mag) | 8.324 ± 0.019 | 2MASS |
| H (mag) | 7.774 ± 0.051 | 2MASS |
| K (mag) | 7.617 ± 0.023 | 2MASS |
| Spectroscopic and Derived Properties | | |
| ξ (mas yr⁻¹) | −343.83 ± 0.073 | Gaia DR2 |
| ξ (mas yr⁻¹) | −573.13 ± 0.073 | Gaia DR2 |
| Barycentric rv (km s⁻¹) | −46.826 ± 0.015 | Gaia DR2 |
| Distance (pc) | 44.583 ± 0.096 | Gaia DR2 |
| Age (Gyr) | 11 ± 2 | This Paper |
| Spectral Type | K3.5V ± 0.5 | This Paper |
| [Fe/H] | −0.47 ± 0.08 | This Paper |
| log g | 4.62 ± 0.02 | This Paper |
| T eff (K) | 4716 ± 60 | This Paper |
| M (M) | 0.888 ± 0.016 | This Paper |
| R (R) | 0.690 ± 0.023 | This Paper |
| L (L) | 0.227 ± 0.003 | This Paper |

We obtained an R ≈ 2000 infrared spectrum of Wolf 503 covering the spectral range between 0.7 and 2.55 μm at the NASA Infrared Telescope Facility (IRTF). We use the SpeX spectrograph in SXD mode with the 0.5 × 15′′ slit. The spectrum was taken on UT 2018 June 03, on a partly cloudy night with an average seeing of 0.6″. Reduction of the spectrum was performed with the SpeXTool (Cushing et al. 2005) and xtellcor (Vacca et al. 2003) software packages as in Dressing et al. (2017). The sky subtraction was performed using a nearby A star, HD 122749, observed immediately after Wolf 503b. The final JHK band IRTF spectra of Wolf 503 are shown in Figure 3 and compared to those of spectral standards. The best visual match indicates a spectral type of K3.5V ± 0.5, suggesting an effective temperature of approximately 4750 ± 100 K from the SpeX spectrum.

During the vetting of candidates from C17 of K2 described in Crossfield et al. (2018), a spectrum was also obtained from the Tillinghast Reflector Echelle Spectrograph (TRES; Füürész 2008) mounted on the 1.5 m Tillinghast Reflector at Fred Lawrence Whipple Observatory on Mount Hopkins on UT 2018 May 23. TRES is a fiber-fed, cross-dispersed echelle spectrograph with a resolving power of R ≈ 44,000, a wavelength coverage of 3850–9100 Å, and radial velocity stability of 10–15 m s⁻¹. The spectrum was reduced and atmospherically corrected according to the procedure described in Buchhave et al. (2010), and we derived stellar atmospheric parameters using the Stellar Parameter Classification code (SPC; Buchhave et al. 2012). We find T eff = 4640 ± 50 K, log g = 4.68 ± 0.10, [Fe/H] = −0.47 ± 0.08, and v sin i = 0.8 ± 0.5. We note that SPC determines the stellar parameters using synthetic spectra with a fixed macroturbulence of 1 km s⁻¹, which may bias v sin i, measurements of slow rotators like this one. Regardless, Wolf
503 has a low projected rotational velocity, as is expected for an old K dwarf, which bolsters its status as a good candidate for precise radial velocity observations. We derive a barycentric radial velocity of $-46.629 \pm 0.075 \, \text{km s}^{-1}$.

We conclude that the SpeX spectrum and the TRES spectrum result in consistent estimates of the stellar temperature. These values are also consistent with the value from the PASTEL catalog of 4759 K (Soubiran et al. 2010), as well as Wolf 503’s colors ($B - V = 1.02, V - K = 2.66$), leading us to adopt the K3.5V±0.5 subtype.

Finally, we adopt $T_{\text{eff}} = 4716 \pm 60$ K, the average and scatter of the three spectroscopic values, as our final value for the stellar temperature. We then calculate the stellar parameters using Isoclassify (Huber et al. 2017). Isoclassify uses measured stellar parameters in comparison to a sample of 2200 Keplar stars with combined Gaia and asteroseismic data in order to determine stellar parameters such as mass and radius with reliable uncertainty based on MIST models. We adopt the log $g$ and [Fe/H] from the TRES spectrum, as well as the $K$ magnitude, which is least affected by extinction. We determine the best stellar radius estimate using the direct method in Isoclassify (Huber et al. 2017), which uses bolometric corrections and direct physical relations to derive stellar properties, but does not return a mass. We obtain the stellar mass using the grid mode, which places the star on stellar evolutionary tracks to determine its properties. The two modes returned consistent stellar radii. The resulting stellar parameters are listed in Table 1.

### 2.4. Target Validation

By far the most pernicious false positives detected in K2 data are eclipsing binaries, which may closely resemble exoplanet transits at grazing incidence, or when the binary system is found in the background of a brighter star (Abdul-Masih et al. 2016). We used archival and adaptive optics images to investigate the possibility of a false positive detection due to a companion star or background sources, and find no source in the vicinity of Wolf 503 that could have contaminated our detection.

#### 2.4.1. Adaptive Optics

Wolf 503 was observed on the night of UT 2018 June 01 UT at Palomar Observatory with the 200'' Hale Telescope using the near-infrared adaptive optics (AO) system P3K and the infrared camera PHARO (Hayward et al. 2001). PHARO has a pixel scale of 0''025 per pixel with a full field of view of approximately 25''. The data were obtained with a narrowband Br-γ filter ($\lambda_0 = 2.18; \Delta \lambda = 0.03 \, \mu m$). The narrowness of the filter enables integration on the primary target without saturation, and the central wavelength of the filter is sufficiently close to the central wavelength of the 2MASS $K_{\text{short}}$ filter ($\lambda_0 = 2.15; \Delta \lambda = 0.31$), enabling the deblending of the 2MASS magnitude of the primary star based on the observed magnitude difference of any detected companions.

The AO data were obtained in a five-point quincunx dither pattern with each dither position separated by 4''. Each dither position is observed three times, each offset from the previous image by 0''5 for a total of 15 frames; the integration time per frame was 4.428 s for a total of 66 s on-source integration time. We use the dithered images to remove sky background and dark current, and then align, flat-field, and stack the individual images. The final PHARO AO data have a FWHM of 0''099. The sensitivities of the final combined AO image were determined by injecting simulated sources azimuthally around Wolf 503 every 45° at separations of integer multiples of the...
FWHM of the central source. The brightness of each injected source was scaled until standard aperture photometry detected it with $5\sigma$ significance. The resulting brightness of the injected sources relative to Wolf 503 set the contrast limits at that injection location. The average $5\sigma$ limits and associated rms dispersion caused by azimuthal asymmetries from residual speckles as a function of distance from the primary target are shown in Figure 4.

The AO imaging revealed no additional stars within 0\textquoteleft 009. For a system at a distance of 44.58 pc, this limits the separation of a possible binary to less than 4.4 au.

### 2.4.2. Archival Images

Even in the absence of a nearby contaminant, adaptive optics cannot eliminate the possibility of a background source directly behind the target, which could be responsible for the signal itself, or could otherwise decrease the apparent transit depth. To address this, we exploit archival imaging from the Palomar Observatory Sky Survey I, and the SERC-EJ and SERC-ER surveys taken on the UK Schmidt telescope. Figure 5 shows the present-day location of Wolf 503 in each of the 3 surveys. The blue plate from POSS I (taken 1952 May 23) and the red SERC-ER survey image (taken on 1993 March 29 with the UK Schmidt Telescope) have a 1\textquotesingle pixel scale, and the blue SERC-EJ image (taken 1983 May 7) has a 0\textquoteleft 059 pixel scale.

The high proper motion of Wolf 503 reveals clearly that there is no background source at the star’s 2018 location. The object detected nearest to Wolf 503’s present-day location is the galaxy LCRS B134447.1-055347, which is located $\approx$25\textquotesingle 1 from the target, placing it outside the aperture used in our extraction. Moreover, the galaxy has a Gaia magnitude of 19.6: being both 10 mag fainter and outside the aperture, we find no background sources that may influence our photometry, indicating that any possible stellar contaminant must be bound within the limit of 4.4 au given by our adaptive optics.

As discussed in Section 2.5, the light curve is consistent either with a transiting planet or a highly specific multiple star system, therefore we find the likelihood of a false positive due to a bound eclipsing binary companion to be extremely low.

One scenario that remains plausible is the case of a bound companion orbiting within 4.4 au that does not transit Wolf 503, but contributes to the total flux and dilutes the planet’s transit depth. According to the distribution of binary star systems found in Raghavan et al. (2010), fewer than 12% of stars belong to such close systems. Additionally, Kraus et al. (2016) found that binary systems with separations smaller than 50 au are not likely to host planets, and that planets in binary systems orbiting closer than 5 au are extremely rare, suggesting that this scenario is also not likely.

Such a companion would also induce a significant radial velocity, of which there is no indication throughout measurements from Gaia, CORAVEL, and our team. Each of these measurements is consistent within 2$\sigma$ and differs by less than 0.8 km s$^-1$. Even a 0.1 $M_\odot$ companion orbiting at 4.4 au would induce a radial velocity of 1.6 km s$^-1$, and according to the modeled mass–luminosity relations in Spada et al. (2013), such a star would be roughly 4 mag dimmer in the $K$ band and 9 mag dimmer in $V$, and would not significantly affect the transit depth. The possibility of such a companion could conclusively be eliminated with high-resolution spectroscopy.

### 2.5. Light-curve Fitting

We fit the light curve of Wolf 503 using ExoFIT, a modular light-curve analysis tool developed for the joint analysis of data from Kepler, Spitzer, and HST. ExoFIT jointly or individually fits transits and explores the parameter space using the Affine Invariant Markov Chain Monte Carlo (AI-MCMC) Ensemble sampler available through the emcee package in Python. Details can be found in Benneke et al. (2017).

We performed individual transit fits in addition to fitting the transits simultaneously. For all fits, we initialize the MCMC chains with uniform priors using the best-fit values from the initial detection pipeline (see Section 2.1), and fit the transit start time $T_0$, duration $T_{14}$, depth $R_p/R_*$, impact parameter $b$, limb-darkening coefficient, and linear background for each transit and scatter term. For the joint fit, we also fit the period $P$. In each fit, we assign 6 walkers for each parameter and find good convergence after 3000 steps, taking the initial 60% as burn-in.
The transits were first fit individually, and the resulting fits are shown in Figure 6. Of the 11 transits observed, all are consistent in \( R_p/R_\star \) and \( T_{14} \). We obtain our best-fitting planet parameters from a joint fit of the 11 transits using the initialization as previously described. The parameters resulting from this fit are summarized in Table 2, where the errors in \( R_p \) and \( a \) are dominated by the stellar parameters. The best-fit light curve is shown in Figure 7, where the combined residuals are well-behaved.

The best fit is distinctly flat-bottomed, inconsistent with the V-shaped light curves characteristic of eclipsing binaries, unless Wolf 503 belongs to a trinary system with two smaller stars orbiting on a 12 day period, within 4.4 au, aligned to be completely eclipsing. In addition to being far more contrived than a single transiting planet, the depth and duration of the transits in Figure 6 are highly regular, and do not show the even–odd variation that would be expected of such an eclipsing binary. As discussed in Section 2.4, such a companion would also induce a significant radial velocity, which has not been detected and would be easily revealed using high-resolution spectroscopy.

### 3. Discussion

From our combined imaging, photometric, and spectral analyses, we establish Wolf 503b as a \( 2.03^{+0.08}_{-0.07} R_\oplus \) planet orbiting its K3.5V±0.5 dwarf host star with a period of 6.0012 days. Wolf 503b is truly distinguished, as its size places it directly at the edge of the radius gap near 1.5–2.0 \( R_\oplus \), while its bright host star (\( H = 7.77 \) mag, \( V = 10.28 \) mag) makes it one of the best targets for radial velocity follow-up and transit spectroscopy at its size (Figure 8).

Radial velocity measurements of Wolf 503b present an excellent opportunity to probe the bulk density of a planet just outside the radius gap. The amplitude of the expected RV signal depends strongly on the planet composition and amount of gas accreted. As Wolf 503b is similar in size to 55 Cnc e, though at a lower temperature, we investigate its composition using the mass–radius relationships for rocky compositions found in Valencia et al. (2010) and Gillon et al. (2012). For the gas-poor scenario, the minimum mass required for a rocky composition (with no iron), is roughly 10 \( M_\oplus \), with an Earth-like composition corresponding to 14 \( M_\oplus \). These masses would result in RV amplitudes of roughly 4.5 and 6.3 m s\(^{-1}\). For a volatile planet with a 0.01% H/He envelope, we would expect a mass of roughly 8 \( M_\oplus \), whereas a 20% water envelope would suggest 6 \( M_\oplus \), and the empirical mass–radius relation by Weiss et al. (2013) would suggest 5.3 \( M_\oplus \), giving RV amplitudes of 3.6, 2.7, and 2.4 m s\(^{-1}\). These amplitudes are detectable with existing precision radial velocity spectrographs, particularly for a bright target such as Wolf 503. As the gas-rich scenario produces much smaller RV amplitudes, these measurements will provide critical constraints on the bulk composition of the planet.

![Figure 6](image-url) Individual K2 transit fits of Wolf 503b. The left panel shows each individual transit with its corresponding best-fit model. The residuals are shown in the center panel, with the residuals in the range \( T_0 \pm T_{14} \) marked in black. The right plot shows the best guess and 1\( \sigma \), or 68% confidence limits on the \( R_p/R_\star \) and \( T_{14} \) parameters, which are consistent for all transits, further supporting the argument that the signal best matches that of a transiting planet. Uncertainties on the first and tenth transits (red and violet) are higher due to masked data points coinciding with a thruster burn near the time of the transit.
Wolf 503b is also an ideal target for detailed characterization with HST and JWST. The signal to noise for future HST transit spectroscopy was estimated in comparison to other confirmed planets near the radius gap, assuming a volatile-rich H/He envelope for each planet. Using the same estimated planet mass of $5.3 M_\oplus$, Wolf 503b is expected to be the second best candidate, behind only 55 Cnc e, for studying a planet in the $1.8-2.1 R_\oplus$ range, where planets may be transitioning into the radius gap through photoevaporation. The planet is also approximately 1000 K cooler than 55 Cnc e, making it more likely to have a significant H$_2$ fraction in its atmosphere, but may also be in the process of photoevaporation. With $J = 8.32$ mag, it is just below the saturation levels of $J > 7$ mag and $J > 6$ mag on the NIRISS and NIRSpec grisms. If Wolf 503b indeed harbors a thick atmosphere, it is one of the best known targets for transmission spectroscopy at its size. Figure 9 shows two simulated transit spectra for Wolf 503b, the blue corresponding to a hydrogen-rich, Neptune-like atmosphere, and the orange corresponding to an atmosphere rich in water. Simulated NIRISS and NIRSpec data for the Neptune-like atmosphere is overplotted, demonstrating the high-confidence with which we will be able to constrain the structure and abundances of atmospheric molecules on Wolf 503b.

Both radial velocity measurements and atmospheric characterization with HST would be valuable short-term follow-up to this work. Wolf 503b is among only a handful of planets in its size range for which this follow-up can be done efficiently today. As such, we expect Wolf 503b to play a critical role in providing near-term insights into the distribution of core masses, the envelope fraction, and the role of photoevaporation for planets near the Fulton gap. It can also serve as an archetype for this class of small planets orbiting nearby stars in preparation for future characterization of similarly bright TESS systems.
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