Ge$_{136}$ type-II clathrate as precursor for the synthesis of metastable germanium polymorphs: A computational study

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The response to compression of the clathrate type-II structure Ge(cF136) is investigated by means of 

I. INTRODUCTION

Group-IVa elements (tetrels) display an extended polymorphism, reflected in a wide range of properties. Due to their technological relevance, semiconducting silicon and germanium have motivated repeated investigations on their polymorphism [1–6]. Germanium displays higher carrier mobility than silicon and finer band-gap tunability [7]. In Ge, lowering of phonon frequencies promotes electron-phonon coupling towards superconductivity [2,3]. Metallization occurs in silicon and germanium upon compression [1]. The possibility of metallic germanium under room conditions is very intriguing and intensively debated [2,8,9], while pressure favors superconductivity in elemental Ge [3].

The polymorphism of germanium echoes in many respects that of silicon [1] by comparatively higher transition pressures [10]. Upon compression, semiconducting Ge(cF8) transforms into the $\beta$-Sn type (tI4, space group I4$_1$/amd) at about 10 GPa [11], followed by an Imma phase [12], a simple hexagonal phase (P6$_3$/mmmc) [13], an orthorhombic Cmca phase [14], and finally, upon further compression above 180 GPa, the hexagonal close-packed arrangement (P6$_3$/mmc) [14]. Additional metastable polymorphs can be expected, if high-pressure experiments are designed to influence nucleation, by changed decomposition protocols, by low temperatures [15], by nonhydrostaticity [16,17], or by choosing a different Ge allotrope as the starting material [18].

Open-framework structures such as clathrates [19–21] offer manifold opportunities to tune their electronic, magnetic, spectral, and transport properties, including thermoelectricity [22–26]; therefore this class of compounds is enjoying renewed interest. The type-II clathrate Ge(cF136) was chemically synthesized from a salt precursor, Na$_{12}$Ge$_{17}$, by mild oxidation with HCl. It is stable under room conditions and persists up to 693 K [27]. Other reported metastable modifications, Ge(tP12) and Ge(cI16) ($\gamma$-silicon type, BC8), can be generated by decomposition [28–30], while Ge(hR8) [17] results from direct compression of cF136. In all these allotropes, the four-bonded atoms adopt nearest-neighbor distances that are similar to those of diamond-type Ge(cF8).

In order to systematically approach the study of phase-transition mechanisms, the use of molecular dynamics accelerated techniques is mandatory to efficiently overcome high energy barriers. Metadynamics [31] explores free-energy surfaces by depositing a history-dependent bias along selected collective variables (CVs). To investigate pressure-induced polymorphism, the whole cell can act as a CV [32–34]. In a previous work [35], we have used metadynamics to investigate structural transformations between diamond Ge(cF8) and Ge(tI4), $\beta$-Sn-type structure. Along the diamond $\rightarrow$ $\beta$-tin transition, the metadynamics visited an intermediate of bct-5 topology (I4/mmm).

This hypothetical structure, in which Ge atoms are five-fold coordinated, is metallic, mechanically stable under room conditions, and superconducting at room temperature. In the metadynamics run, only one box parameter was markedly affected, suggesting nonhydrostatic shearing as the protocol of choice when attempting to obtain bct-5.

Despite the noticeable differences in atomic volumes (24.5 and 28.3 Å$^3$, respectively), the compressibility of Ge(cF136), as revealed by the bulk moduli $B_0 = 76 \pm 6$ GPa [17], is very close to the value of cubic diamond Ge(cF8), $B_0 = 75$ GPa.
FIG. 1. Enthalpy differences between Ge\(_{136}\) [Ge(cF136)], \(\beta\)-tin [Ge(tI4)], and bct-5. Ge\(_{136}\) represents the baseline. The pressure values considered in this paper are indicated as vertical lines at 2.5 and 5.0 GPa.

[36]. Total energy calculations (using the density functional method SIESTA [37]) indicate that Ge(cF136) is the energetically lowest-lying modification among metastable germanium polymorphs.

The transition pressures between cF136, cI4, and bct-5 were evaluated based on the enthalpy equivalence between phases, \(E_1 + pV_1 = E_2 + pV_2\). The difficulty in accessing bct-5 from direct, isotropic compression lies in the presence and higher relative stability of \(\beta\)-tin. A strategy must therefore be designed that takes into account nonhydrostatic compression (see above and Ref. [35]) of a suitable intermediate as learned from a recent study [38]. In the latter, the pressure-induced amorphization of Ge\(_{136}\) and subsequent fast recrystallization into \(\beta\)-tin or transformation into low-density amorphous (LDA) Ge were investigated using \emph{ab initio} molecular dynamics and metadynamics. Metadynamics was used up to the occurrence of amorphous intermediates in the 136-atom box, followed by plain equilibrium molecular dynamics protocols until phase crystallization. Fast recrystallization was also observed in very high density amorphous (VHDA) Si [39] recently, likely owing to common tetrahedral structural motifs.

In this paper, differently from Ref. [38], we apply metadynamics beyond the stage of system amorphization to systematically identify configurations kinetically accessible from disordered intermediates. Two simulation pressures were therefore selected, \(p = 5.0\) GPa and \(p = 2.5\) GPa, above and below the threshold pressure for \(\beta\)-tin [Ge(tI4); Fig. 1]. All small-cell metadynamics runs were performed based on the 34-atom primitive cell of Ge\(_{136}\). The choice of a small simulation box shall favor anisotropic box fluctuations and also assist disordered-ordered system transitions as demonstrated by small-cell molecular dynamics applied to ice polymorphism [40].

II. RESULTS

In Fig. 2, five metatrajectory snapshots from type-II clathrate Ge(cF136) to Ge(tI4) at \(p = 5.0\) GPa and \(T = 300\) K are shown, while Fig. 3 summarizes the evaluated descriptors for the same metatrajectory. The percentage of pure \(\beta\)-tin species is displayed on the opposite y axis in Fig. 3 (indigo curve). The metasteps corresponding to the snapshots of Fig. 2 are indicated in Fig. 3 as thick marks on the horizontal axis and labeled (Ge\(_{136}\), I1, I2, I3, and \(\beta\)-tin).

The stability of Ge(cF136) is reflected in the large number of metasteps spent lingering in the initial basin. Only after more than 100 ps (about 210 metasteps) a sudden enthalpy drop (Fig. 3, main panel) is observed, as the system moves into a regime of amorphous but rather dense structures [Fig. 2(b) and intermediate I1 in Fig. 3]. The average circuit size (ACS) [41] and enthalpy display, in fact, minima, while the coordination sequence (CS; Fig. 3, right inset), which is a set of integers \(\{n_1, n_2, \ldots\}\), where \(n_i\) indicate the number of atoms separated from a central one by a minimal path of \(i\) bonds, rapidly increases. Subsequently, the system remains of mixed
FIG. 4. (a)–(e) Snapshots from a metadynamics trajectory at \( p = 2.5 \text{ GPa} \) and \( T = 300 \text{ K} \). The corresponding time steps are indicated in Fig. 5 as thick marks on the horizontal axis. \( \text{Ge}_{136} \) (a) transforms into \( \text{bct-5} \) (e). An unstable clathrate structure \( \text{Ge(mC68)} \) is visited (b). In (a) and (b), selected Ge atoms are highlighted in green to stress the closeness between the two structural motifs. The characteristic square-pyramidal coordination polyhedron of \( \text{bct-5} \) is shaded in green in (e).

structural character, as measured by the above indicators, with a single attempt at forming \( \beta \)-tin (indigo peaks) around metastep 265.

The ACS peak between metasteps 295 and 305 corresponds to a marked drop in the coordination sequence. In this region, thermal fluctuations bring the system close to \( \beta \)-tin (Fig. 3, \( \beta \)-tin %); however, the structure remains rather expanded and does not fully lock in. The configuration at metastep 332 [Fig. 2(c) and intermediate I2 in Fig. 3] is representative for the subsequent 40 metasteps and contains mixed structural motifs, comprising remaining \( \beta \)-tin-like features (4+2 coordination) and higher coordination numbers up to 12.

Between metasteps 370 and 400, plateaus can be recognized by inspection of the ACS, CS, and enthalpy graphs. This region features structures that are related to \( \alpha \)-polonium type (cP1) but distorted (while Ge is on average sixfold coordinated, the angles of the octahedral coordination polyhedron deviate from 90°) and more dense, as displayed by the coordination sequence average 6, 18, 42 against the expected 6, 18, 38 for ideal \( \alpha \)-polonium (cP1). A similar occurrence of a denser structural motif is observed between steps 405 and 425, where a simple hexagonal structural pattern is closely matched [Fig. 2(d) and intermediate I3 in Fig. 3], with a slightly expanded volume though, as can be inferred from the coordination sequence 7, 5, 25, 54 against 8, 26, 56 of ideal hP1. Cell fluctuation enhancements are characteristic of the metadynamics methods and contribute to driving the systems to the final \( \beta \)-tin basin, into which the whole system has transformed [Fig. 3, \( \beta \)-tin %, and Fig. 2(e)] after about 500 metasteps.

The same descriptors were applied to a metatrajectory at \( p = 2.5 \text{ GPa} \) and \( T = 300 \text{ K} \) that connects \( \text{Ge(cF136)} \) to \( \text{bct-5} \) (Figs. 4 and 5). The percentage of pure \( \text{bct-5} \) appears on the opposite y axis in Fig. 5 (blue curve). The lower-pressure regime markedly affects the metatrajectory. After about 175 metasteps, \( \text{Ge(cF136)} \) distorts into an intermediate, which can be quenched [Fig. 4(b) and intermediate I1 in Fig. 5].

We characterized this intermediate with the Pearson symbol \( \text{Ge(mC68)} \). This clathratelike structure is higher in energy than \( \text{Ge(cF136)} \) \( \Delta E = E(\text{cF8}) - E(\text{cF136}) = 0.03 \text{ eV/atom} \), while \( \Delta E = E(\text{cF8}) - E(\text{mC68}) = 0.16 \text{ eV/atom} \) and can be relaxed in all parameters at 0 K. The chosen projection of Figs. 4(a) and 4(b) allows us to appreciate the structural closeness of \( \text{Ge(cF136)} \) and \( \text{Ge(mC68)} \). A set of atoms has been highlighted [green spheres in Figs. 4(a) and 4(b)] to illustrate the formation of chains as part of the transforma-
tion mechanism, which leaves large portions of Ge(cF136) unaffected. A short molecular dynamics run started from Ge(mC68) \((t = 10 \text{ ps}, p = 0 \text{ GPa}, T = 300 \text{ K})\) did not show any appreciable structural deformation. Nonetheless, imaginary frequencies were found in the phonon dispersion analysis (not shown), suggesting overall mechanical instability. In fact, after 25 metasteps, mC68 evolves into another structural area. Here a single structural motif can hardly be extracted, for what can be described as intermediate configurations between Ge(cF136) and a denser, \(\beta\)-tin-like phase: The average circuit size for this segment is \(\sim 4.9\), intermediate between 5.1 for Ge(cF136) and 4.67 for \(\beta\)-tin; the coordination sequence has also increased with respect to Ge(cF136), indicating overall denser structures [Fig. 4(c) and intermediate I2 in Fig. 5], while the enthalpy also takes intermediate values.

Thereafter, a long segment of the trajectory of about 175 metasteps consists of amorphous structures that occasionally show local order in correspondence to maxima in ACS and in CS, and to enthalpy minima. A representative snapshot is displayed in Fig. 4(d) (Fig. 5, intermediate I3).

At metastep 430 the metadynamics trajectory reaches the basin of bct-5, as indicated by blue peaks in Fig. 5. Its characteristic fivefold coordination is illustrated by a green pyramidal polyhedron in Fig. 4(e). bct-5 appears in a region of higher enthalpy values, corresponding to a tiny local minimum. The bct-5 motif forms within a landscape dominated by low-crystalline motifs, from which \(\beta\)-tin crystallization is not observed.

Two relevant pressure protocols can therefore be identified. At relatively low pressure \((P = 2.5 \text{ GPa})\), a transition from Ge\(_{136}\) to bct-5 can be observed, proceeding over amorphous intermediate steps. While higher pressure always favors the formation of Ge(tI4) from Ge(cF136), lower pressures admit bct-5 among the accessible configurations. Nonetheless, even in a pressure regime below its stability range such as the one implemented here, Ge(tI4) remains a common transient motif linking fourfold to higher coordinations. This shows the need for a thorough investigation of the factors that may determine the appearance of bct-5.

The metadynamics from Ge(cF136) was repeated at 5.0 GPa and 77 K, using the same protocol as for the previous runs. Lowering temperature caused the \(\beta\)-tin phase to sharply lock in, with a clear signature on the enthalpy profile, and this phase remained well defined over many metasteps, as shown in Fig. 6. Also, the transition was markedly smoother as indicated by the gentle drop in the enthalpy profile, the overall progressive drop in the ACS, and the increase in the CS.

As the system escaping the Ge(cF136) well, it visits an open-framework structure (metasteps 300–310) which is denser than the initial structure and which displays a minimum in the ACS and an increased CS value (Fig. 6, left and right insets). This is less of a metastable phase than an activated intermediate, which is followed by a region of amorphous structures, as any coherent structural motif is hardly distinguishable, while neither ACS nor CS indicate the existence of a clearly defined intermediate crystalline structure.

Around metastep 380 the system recrystallizes into a \(\beta\)-tin-like structure (maximum in ACS and minima in CS values and enthalpy). Cell fluctuations bring it back to an amorphous stage, which is the gateway to the exact \(\beta\)-tin basin. Between two 100% \(\beta\)-tin peaks centered at metasteps 416 and 480, other amorphous phases are touched upon. Nonetheless, it can be expected that these effects can be further affected by retuning the choice of the Gaussian parameters in the simulation setup.

The metadynamics run at \(p = 2.5 \text{ GPa} \) and \(T = 77 \text{ K}\) was performed with a slightly modified protocol, which entailed rescaling the width and height of the Gaussian bias (see Sec. V for details). The rescaling to a lower value would, namely, allow for a longer persistence in stable basins. A full analysis based on the same descriptors as for the other runs (enthalpy, ACS, and CS), augmented by the statistics of fivefold rings (green curve), is presented in Fig. 7, main panel. Fivefold rings introduce an additional means to sharply distinguish between Ge(cF136), Ge(tI4), and bct-5. Their characteristic number of fivefold rings is 180 and 136, respectively (orange labels on the opposite \(y\) axis in the main panel of Fig. 7), while there are no fivefold rings in bct-5.

In Ge(cF136) fivefold rings occur as so-called \(K_5\) graphs [42], in which six fivefold rings share each two edges around a Ge atom. They constitute therefore characteristic building blocks based on fivefold-ring statistics, whose monitoring is therefore an excellent indicator of structure evolution.

Considering major changes only, three regions can roughly be distinguished by enthalpy in Fig. 7: A first segment from metastep 1 to metastep 660, dominated by Ge(cF136), a second region between metasteps 660 and 1330, and a third region (metasteps 1580–3260) of overall lower enthalpy (including fluctuations). The CS indicates that the first two regions are dominated by tetrahedral motifs (first CS = 4, Fig. 7, right inset, black curve), while denser structures are found in the third region. Therein, the number of fivefold rings remains largely within (occasionally above) a well-defined region (orange dashed lines in Fig. 7, main panel), while depletion of that figure occurs only episodically. In Fig. 8, eight snapshots from the metatrajectory are presented. Their corresponding metasteps are indicated on the main horizontal
FIG. 7. Evolution of average circuit size (left inset), coordination sequence (right inset), enthalpy per atom (black curve, main panel), and number of fivefold rings (green curve, main panel), for a metadynamics run at $p = 2.5$ GPa and $T = 77$ K. Eight thick marks on the horizontal axis of the main panel indicate the metasteps relative to the eight snapshots of Figs. 8(a)–8(h). On the opposite y axis in the main panel, the characteristic number of fivefold rings for Ge$_{136}$ and β-tin is indicated: 180 and 136, respectively (orange labels and dashed orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines). In bct-5, there are no fivefold rings. The progress of the conversion from Ge$_{136}$ to bct-5 is indicated as a blue (turquoise) orange lines).
FIG. 8. (a)–(h) Snapshots from a metadynamics trajectory at $p = 2.5$ GPa and $T = 77$ K. The corresponding time steps are indicated in Fig. 7 as thick marks on the horizontal axis in the main panel. Ge$_{136}$ (a) transforms into bct-5 (g) over several intermediates [(b)–(f)]. The region of stability of Ge$_{136}$ (a) is followed by a region of substantially amorphous motifs of intermediate enthalpy values (b), which in turn evolves into a regime of lower enthalpy on average [(c)–(h)]. The latter is characterized by the occurrence of local fivefold coordination. Therein, $\beta$-tin is occasionally visited but fails to lock in at any point.

phases, but as transient motifs, particularly around bct-5, to which it can be mechanistically related in this low-pressure, low-temperature regime.

Based on our detailed mechanistic analysis, and previous results on the role of amorphization [38] at higher pressure, a synthetic pathway for bct-5 would therefore require separate steps, including (i) the amorphization of the clathrate motif at high pressure followed by (ii) recrystallization at a lower pressure such as 2.5 GPa. A possible way to accomplish this protocol might be a short shock wave, which would destabilize the clathrate framework, followed by a low-pressure regime, allowing the system to recrystallize. Alternatively, the application of anisotropic stress to the clathrate could induce both the amorphization at lower pressure and the crystallization to bct-5.

IV. CONCLUSIONS

The compression of Ge$_{136}$ [Ge(cF136)] was studied by means of ab initio small-cell metadynamics. The choice of this approach, which makes no assumption regarding the final state, was meant to investigate the feasibility of recrystallization of the elusive bct-5 phase from specific precursors and to study under what conditions this pathway could be preferred. At $p = 5.0$ GPa and $T = 300$ K, the preferential product was identified as $\beta$-tin [Ge(tI4)]. At lower pressure ($p = 2.5$ GPa and $T = 300$ K), the fivefold-coordinated bct-5 appeared following pronouncedly anisotropic cell fluctuation. To shed further light on the competition between these two structural motifs, metadynamics runs were also performed at lower temperature, $T = 77$ K. $\beta$-tin remained accessible through direct compression at higher pressures, while bct-5 appeared exclusively at lower pressure values. Additionally, lower-temperature metadynamics discloses a rich landscape of different, transient fivefold-coordinated structural motifs, which can in principle crystallize from amorphous precursors, generated by direct compression of Ge$_{136}$. Unlike $\beta$-tin, bct-5 does not contain any fivefold rings (see Fig. 8), which correlates with its vibrational entropy [45], and has a comparatively higher enthalpy per atom. The combination of low pressure, low temperature, nonhydrostatic compression, and sufficient time to allow for crystallization can therefore be explored as a gateway for the synthesis of this elusive phase, while higher pressures appear to consistently favor $\beta$-tin [Ge(tI4)] [38] due to fast crystallization kinetics.

V. METHODS

A. Metadynamics

Metadynamics [31–34] allows for the exploration of the energy surface along one or more collective reaction coordinates. The method is independent of the level of theory used, it does not require prior knowledge of the energy landscape, and its sampling efficiency can be enhanced by parallel runs started from different configurations. The time evolution of the system is biased by a history-dependent potential, which discourages the system from visiting already harvested
regions of the potential [46]. Efficiency is achieved in metadynamics also through dimensionality reduction. Instead of studying the problem in the full 3N-dimensional configuration space of N particles, a relatively small number of collective coordinates $s = (s_1 \cdots s_m)$ are used instead, which provide a coarse-grained description of the system and are able to distinguish between different free-energy minima, i.e., different phases. The inclusion of slow degrees of freedom in the space of collective variables positively impacts the performance of the method.

Each metadynamics metastep consisted of a molecular dynamics run in the NVT ensemble for a total simulation time of 0.5 ps (time step 2 fs) at either 300 K or at liquid nitrogen temperature, 77 K, depending on the regime studied. All metadynamics runs were performed with 34 atoms in the simulation box, which served as a collective (six-dimensional) variable. The size of the minimal box ensured commensurability of all already known phases, either open or dense, including Ge(cF136).

B. Density functional computational layers

SIESTA [37] was used as the density functional theory (DFT) molecular dynamics layer. For all compression protocols, electronic states were expanded by a single-$\zeta$ basis set constituted of numerical orbitals with a norm-conserving Troullier-Martins pseudopotential description of the core electrons [47]. The single-$\zeta$ basis set dramatically reduces computational times providing, nonetheless, the right topol-electronic part.

To accurately trace the structure evolution in the metadynamics simulations, the average size of the shortest circuits in the structures was calculated [48]. The average circuit size for Ge(cF136) is 5.12, while for $\beta$-tin it is 4.67. The average circuit size should decrease along the transition from the clathrate phase to the $\beta$-tin phase due to a reduction in the number of fivefold rings.

To discriminate between $\beta$-tin and bct-5 networks, coordination sequences (CSs, up to the third shell) were calculated at each metastep. For $\beta$-tin the CS reads (one atom type) as follows: 6, 22, 46. For bct-5 it is (one atom type) 5, 16, 33. Coordination sequences were also used to estimate the percentage of a phase (either $\beta$-tin or bct-5) along a metadynamics run, by looking for atoms with either $CS_{bct-5}$ 5, 16, 33 or $CS_{\beta-tin}$ 6, 22, 46.

In the case of new structures, ideal space group and asymmetric units were identified with the Generation, Analysis and Visualization of Reticular Ornaments using GAVROG Symmetry, Structure (Recognition) and Refinement (GAVROG SYSTRE) package [49].

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Herein, we use the term “shortest circuits” to emphasize that they are in general different from rings, which are often used to characterize the local structure of amorphous materials. In short, rings are circuits with the special property that any circuit can be written as a disjunctive sum of rings, i.e., rings form a basis for the cycle space of a graph.