The domain and property illusion of anomalous anisotropic electric conductivity

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Abstract

The unique determination of electrical conductivity is extensively studied for isotropic conductivity ever since Calderon’s suggestion of the EIT (Electrical Impedance Tomography) problem. However, it is known that there are many anisotropic conductivities producing the same Dirichlet-to-Neumann map; moreover the anisotropic conductivities giving the same Dirichlet-to-Neumann map are classified using the equivalence relation with respect to the change of variables. The change of variable argument is applied to the theory of near-cloaking.

We are under an illusion that the domain of anomaly is of a much

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smaller size than actually it is. For this paper, we considered not only
the illusion of the domain of the anomaly, but also the illusion of the
property of the anomaly when the background anisotropic conductivity
is known.

1 Introduction

Recently, the idea of physical devices related to cloaking or invisibility is
suggested more concretely than before and draws much attentions from re-
searchers in these areas \[27, 28, 32\]. Invisibility and cloaking are closely
related with the nonuniqueness of coefficients of the inverse problem mod-
elled by EIT \[3, 8, 14, 15, 21\], acoustic scattering \[2, 20, 31\], electromagnetic
scattering \[12, 24\], and quantum scattering \[11, 13\]. In this paper, only the
inverse problem modelled by EIT is considered.

Voltage potential in a Lipschitz domain \(\Omega\) for an electrical current on
the boundary \(\partial \Omega\) is given as follows:

\[-\nabla \cdot (\sigma \nabla u) = 0 \quad \text{in } \Omega, \quad (1.1a)\]

\[u = f \quad \text{on } \partial \Omega. \quad (1.1b)\]

If \(f \in H^{1/2}(\partial \Omega)\) and

\[0 < L < \frac{y^t \sigma(x)y}{y^t y} < U \quad (1.2)\]

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for all $x \in \Omega$, $y \in \mathbb{R}^n \setminus \{0\}$ and some constants $L$ and $U$, it is known to be $u \in H^1(\Omega)$ [10]. Therefore, we can define the Dirichlet-to-Neumann map $\Lambda_\sigma : H^{1/2}(\partial \Omega) \to H^{-1/2}(\partial \Omega)$ as $\Lambda_\sigma(f) = \nu \cdot (\sigma \nabla u)|_{\partial \Omega}$ using the boundary trace operator, when $\Omega$ is a Lipschitz domain.

EIT is formulated to find $\sigma$ such that $\Lambda_\sigma = \Lambda$, given the Dirichlet-to-Neumann map $\Lambda$. In finite measurements case, EIT is to find $\sigma$ such that $\Lambda_\sigma(f_i) = \Lambda(f_i), i = 1, \ldots, N$ for given Dirichlet and Neumann boundary measurement pairs $(f_i, \Lambda(f_i))_{i=1,\ldots,N}$.

Suppose that $D$ is a Lipschitz domain compactly embedded in $\Omega$ such that $\Omega \setminus \overline{D}$ is connected. Reminding all the conductivities in Table 1 satisfies (1.2), we classify conductivity $\sigma$ in (1.1) into six cases in Table 2: Four cases of them (Case 1, 2, 4, and 5) includes the unknown obstacle $D$ inside the domain of interest $\Omega$, which have a different conductivity from the background $\Omega \setminus \overline{D}$. In other words, $\sigma$ is discontinuous along $\partial D$ in these four cases.

We explain briefly the results on the uniqueness and nonuniqueness in Section 2. In Section 3, several results are given not only for illusion and cloaking of domain $D$ but also for property $B(x)$ illusion.
Table 1: Electrical conductivities used in background and target obstacles

| Case   | Formula                                      |
|--------|----------------------------------------------|
| 1      | $\sigma = a + (b - a)\chi_D$               |
| 2      | $\sigma = a(x) + (b(x) - a(x))\chi_D$      |
| 3      | $\sigma = b(x)$                             |
| 4      | $\sigma = A + (B - A)\chi_D$               |
| 5      | $\sigma = A(x) + (B(x) - A(x))\chi_D$      |
| 6      | $\sigma = B(x)$                             |

Table 2: Electrical conductivities to be reconstructed on the whole domain of interest $\Omega$
2 Uniqueness and nonuniqueness

2.1 Uniqueness

Note that in the six cases in Table 2, the unknowns are $b, b(x), B, B(x)$ and $D$, whereas $a, a(x), A$ and $A(x)$ are assumed to be known. The uniqueness studies of EIT in these six cases proceeded in the following ways:

- Case 1 ($\sigma = a + (b - a)\chi_D$):

  One or two measurements uniqueness results are known when $b$ is assumed to be known [6, 9, 18]. The finite measurement uniqueness is known only for Case 1. In this case, the number of measurements, the geometry of the obstacle $D$, and the choice of suitable Dirichlet or Neumann data minimizing the number of measurements are interesting issues.

- Case 2 ($\sigma = a(x) + (b(x) - a(x))\chi_D$) and Case 4 ($\sigma = A + (B - A)\chi_D$):

  In [17], an orthogonality relation between the two solutions of (1.1) for arbitrary obstacles $D_1$ and $D_2$ is derived. Based on this relation and the Hahn-Banach theorem, a modified orthogonal relation is derived and the uniqueness for Case 2 and Case 4 is proved. In Case 4, additional condition such that $B - A$ is positive-definite is used: this additional condition is generalized and removed by [16] and [23]. The orthogonal
and modified orthogonal equation is also used in Case 5.

- Case 3 \( (\sigma = b(x)) \):

This is a generalization of Case 1 and Case 2, since we do not need to have the known background conductivity \( a(x) \). Many mathematicians conduct extensive works in this case and thus, we are able to understand the unique determination of electrical conductivity when \( b(x) \in L^\infty(\Omega) \) for the two-dimension [4, 7, 30] and \( b(x) \in C^{3/2}(\Omega) \) for three and higher dimensions [1, 29, 35, 36].

- Case 5 \( (\sigma = A(x) + (B(x) - A(x))\chi_D) \):

The uniqueness of \( D \) is shown when \( B(x) - A(x) \) is positive-definite [17]; it is also presented when the positive-definite condition is removed in [22] and the nonuniqueness of \( B(x) \) is pointed out by using the nonuniqueness result of Case 6

- Case 6 \( (\sigma = B(x)) \):

\( B(x) \) is not unique for the given Dirichlet-to-Neumann map. The number of \( B(x) \) making the same Dirichlet-to-Neumann map depends on the number of the diffeomorphism on \( \Omega \), which will be described in Definition 1.
2.2 Nonuniqueness in Case 6

The nonuniqueness of EIT is observed early in [19]. Before stating the result, let us define the diffeomorphism on Ω as follows:

**Definition 1** A map \( F \) defined almost everywhere from \( Ω \) into \( Ω \) itself is called diffeomorphism on \( Ω \), if the following conditions are satisfied:

1. \( F \) is bijective \( \quad (2.3a) \)
2. \( F \) is an identity on \( ∂Ω \) \( \quad (2.3b) \)
3. Jacobian \( DF \) is defined and bounded \( \quad (2.3c) \)
4. \( \det(DF) \neq 0 \) \( \quad (2.3d) \)

**Proposition 2** Let \( F \) be a diffeomorphism on \( Ω \). Then, we have

\[
Λ_{σ} = Λ_{F∗σ},
\]

where

\[
F∗σ = \left(\frac{(DF)σ(DF)^t}{\det(DF)}\right).
\]

**Proof** Knowing the Dirichlet-to-Neumann map is equivalent to knowing the quadratic form

\[
Q(v) = \int σ \nabla v \cdot \nabla v \quad v ∈ H^1(Ω)
\]
by the polarization identity. Then, the proposition follows from the change of variable method for the quadratic form

$$
\int_{\Omega} \sum_{ij} \sigma_{ij} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \, dx = \int_{\Omega} \sum_{ij} \sigma_{ij} \frac{\partial u}{\partial y_k} \frac{\partial u}{\partial y_l} \frac{\partial y_k}{\partial x_i} \frac{\partial y_l}{\partial x_j} \det \left( \frac{\partial x}{\partial y} \right) \, dy \tag{2.5}
$$

as explained in [21].

$F_*$ in Proposition 2 is called the 'push forward map' for a diffeomorphism $F$ on $\Omega$. After Proposition 1 is known, the following questions in Case 5 and Case 6 are raised:

- Is the change of variable (2.4) a unique obstruction to the uniqueness for Case 6? In other words, does $\Lambda_{\sigma_1} = \Lambda_{\sigma_2}$ imply $\sigma_2 = F_* \sigma_1$ for some diffeomorphism $F$ on $\Omega$?

The answer is yes in Case 6, which means that there are no nonuniqueness cases other than the change of variables. Many researchers, including Gunter Uhlmann, proved this problem for the following subcases:

- $n = 2$ and $\sigma \in C^3(\Omega)$ [34]
- $n = 2$ and $\sigma \in C^{0,1}(\Omega)$ (Lipschitz functions) [33]
- $n = 2$ and $\sigma \in L^\infty(\Omega)$ [5]
- $n \geq 3 : \partial D$ and $B$ are analytic [25, 26]

The answer for Case 5 is given in the following section.
3 Illusions of domain ($D$) and property ($B$) in Case 5 and Case 6

In Case 5 ($\sigma = A(x) + (B(x) - A(x))\chi_D$), the uniqueness of $D$ is solved in [16, 17, 22]; yet, the nonuniqueness of $B(x)$ inside $D$ is only commented in [22] using the result for Case 6 (Section 2.2). Theorem 1 in [22] is rephrased in the following theorem:

**Theorem 3** Assume that

$$A(x), A^{1/2}(x), B_i(x), B_i^{1/2}(x)$$
are Lipschitz continuous on $\Omega$ \hspace{0.3cm} (3.6)

and there is a jump on $\partial D$ such that

$$\det(B_i) \neq \det(A) \text{ for } n = 2 \text{ and } B_i \neq A \text{ for } n \geq 3$$ \hspace{0.3cm} (3.7)

for $i = 1, 2$. Then, $\Lambda_{\sigma_1} = \Lambda_{\sigma_2}$ implies

$$D_1 = D_2.$$ \hspace{0.3cm} (3.8)

Denoting $D = D_1 = D_2$, on the boundary $\partial D$, we have

$$\det(B_1) = \det(B_2) \text{ for } n = 2, \text{ and } B_1 = B_2 \text{ for } n \geq 3.$$ \hspace{0.3cm} (3.9)

If we further assume that

$$B_1 = B_2 \text{ on } \partial D \text{ for } n = 2,$$ \hspace{0.3cm} (3.10)

$\partial D$ is connected and analytic and $B|_D$ is analytic for $n \geq 3$. 9
Then, there is a diffeomorphism \( F \) on \( D \) such that

\[
B_2 = F_* B_1. \tag{3.11}
\]

From Theorem 3, in Case 5, the domain \( D \) is uniquely determined by the Dirchlet-to-Neumann map; however, the property \( B \) is nonunique for the same Dirichlet-to-Neumann map. Specifically, the domain of anomaly cannot be cloaked into a much smaller domain, but the property could be illued into other property. In other words, in Case 5, the size illusion is impossible but the property illusion is possible.

### 3.1 Illusion of domain \( D \) for Case 6

Consider the conductivity

\[
\sigma = A(x) + (B(x) - A(x))\chi_D.
\]

In Case 5, \( A(x) \) is the known property. Then, by Theorem 3, there is no cloaking or illusion for domain \( D \). However, if \( A(x) \) is assumed to be unknown, \( \sigma \) is a special case of Case 6 with \( \sigma \in L^\infty(\Omega) \). In this case, there is illusion or cloaking of domain \( D \). Specifically, there is a diffeomorphism \( F \) on \( \Omega \) which illude the domain \( D \) into arbitrary small domain \( U_\epsilon \) through the following conductivity:

\[
\sigma_F(x) = 1 + (F_* B(x) - 1)\chi_{U_\epsilon}(x)
\]

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For example, if $\Omega, D$ and $U$ are balls of radius 2, 1, and $\epsilon > 0$, respectively, centered at $O$, we could take the diffeomorphism $F_\epsilon$ on $\Omega$ as the inverse of the followings:

$$F_\epsilon(x) = \epsilon x \chi_D(x) + ((2 - \epsilon)|x| - (2 - 2\epsilon)) \frac{x}{|x|} \chi_{\Omega \setminus \overline{D}}(x)$$

(3.12)

where $F_\epsilon$ maps $D$ into $U$. If $\epsilon$ goes to 0, we stae that $D$ is ”perfectly cloaked” by the corresponding diffeomorphism $F_0$.

Let us interpret the cloaking in a different point of view. Let $F$ be a diffeomorphism on $\Omega$ such that $F(D_1) = D_2$. Define three kinds of conductivities as follows:

$$\sigma_1 = I_n \chi_{\Omega \setminus D_1} + B_1 \chi_{D_1},$$

(3.13a)

$$\sigma_2 = I_n \chi_{\Omega \setminus D_2} + B_2 \chi_{D_2},$$

(3.13b)

$$\sigma_3 = A_3 \chi_{\Omega \setminus D_3} + B_3 \chi_{D_3},$$

(3.13c)

where $I_n$ is a $n$– dimensional identity matrix.

**Theorem 4** Suppose that

$$D_1 \neq D_2$$

(3.14)

and (3.6), (3.7) hold. Then

$$\Lambda_{\sigma_1} \neq \Lambda_{\sigma_2}$$

(3.15)

Suppose that

$$A_3 = F_* 1, \ B_3 = F_* B_1, \ \text{and} \ D_3 = F_* D_1$$

(3.16)
for some diffeomorphism $F$ on $\Omega$. Then

$$\Lambda_{\sigma_1} = \Lambda_{\sigma_3}$$

(3.17)

Conversely, if (3.17) holds in the two dimension, then there is diffeomorphism $F$ on $\Omega$ satisfying (3.16).

**Proof** (3.15) and (3.17) are followed by Theorem 3 and Proposition 2. The converse is followed by [5].

From (3.15), we do not have any illusion of domain $D_1$ to another domain $D_2$ in Case 5, when the background conductivity $A$ is known. However, if we do not know the background conductivity $A$ in Case 6, we could illude domain $D_1$ into a more smaller domain $D_3$ (or perfectly cloak $D_1$) by the diffeomorphism $F$ on $\Omega$, for example $F_\varepsilon$ in (3.12), using (3.17) in Theorem 4.

### 3.2 Illusion of property $B$ for Case 5

Although domain ($D$) illusion is impossible in Case 5, property ($B$) illusion is possible by the following theorem:

**Theorem 5** Let

$$\sigma_i = I_n + (B_i - I_n)\chi_D, \ i = 1, 2.$$  

(3.18)
If

\[ B_2 = F_* B_1 \quad (3.19) \]

for some diffeomorphism \( F \) on \( \Omega \). Then,

\[ \Lambda \sigma_1 = \Lambda \sigma_2. \quad (3.20) \]

Conversely, if (3.20) holds, then there is diffeomorphism \( F \) on \( D \) such that (3.19) holds under the conditions (3.6),(3.7),and (3.10).

**Proof** (3.20) follows Proposition 2. And the existence of the diffeomorphism \( F \) on \( D \) follows (3.11) in Theorem 3.

From Theorem 5, property illusion is possible up to the change of variables by diffeomorphisms on \( \Omega \). In [24], a few interesting property illusions are considered; The optical transformation of an object into another object is considered. In the paper, stereoscopic image of a man is transformed into an illusion image of a woman and dielectric spoon of the electric permeability 2 into an illusion image of metallic cup of electric permeability -1 in an electromagnetic scattering problem.

### 4 Conclusions

In EIT, the uniqueness for isotropic conductivities is summarized. It is shown that the nonuniqueness for anisotropic conductivities is possible if
and only if the conductivity is pulled back by the diffeomorphism on the domain of interest. In Case 6, when the background is also unknown, illusion and perfect cloaking of domain is possible. In Case 5, when the background is known, illusion of domain is impossible; yet, the illusion of property inside the domain is possible by the pull-back of the diffeomorphism on the domain.

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References

[1] G. Alessandrini. Stable determination of conductivity by boundary measurements. Appl. Anal., 27:153–172, 1988.

[2] H. Ammari, H. Kang, H. Lee, and M. Lim. Enhancement of near cloaking part II: The Helmholtz equation. Comm. Math. Phys. To appear.

[3] H. Ammari, H. Kang, H. Lee, and M. Lim. Enhancement of near cloaking using generalized polarization tensors vanishing structures, part I: The conductivity problem. Comm. Math. Phys. To appear.
[4] K. Astala and L. Päivärinta. Calderon’s inverse conductivity problem in the plane. *Ann. Math.*, 163:265–299, 2006.

[5] K. Astala, L. Päivärinta, and M. Lassas. Calderon’s inverse problem for anisotropic conductivity in the plane. *Commun. PDE*, 30:207–224, 2005.

[6] B. Barcelo, E. Fabes, and J. K. Seo. The inverse conductivity problem with one measurement: uniqueness for convex polyhedra. *Proc. AMS*, 122:183–189, 1994.

[7] R. M. Brown and G. Uhlmann. Uniqueness in the inverse conductivity problem for non smooth conductivities in two dimensions. *Commun. Partial Diff. Eq.*, 22:1009–1027, 1997.

[8] K. Bryan and T. Leise. Impedance imaging, inverse problems, and Harry Potter’s cloak. *SIAM Rev.*, 52:359–377, 2010.

[9] A. Friedman and V. Isakov. On the uniqueness in the inverse conductivity problem with one measurement. *Ind. Univ. Math. J.*, 38:553–580, 1989.

[10] D. Gilbarg and N. S. Trudinger. *Elliptic partial differential equations of second order*. Springer-Verlag, 1983.
[11] A. Greenleaf, Y. Kurylev, M. Lassas, and G. Uhlmann. Schrödinger’s hat: electromagnetic, acoustic and quantum amplifiers via transformation optics. To appear.

[12] A. Greenleaf, Y. Kurylev, M. Lassas, and G. Uhlmann. Cloaking devices, electromagnetics wormholes, and transformation optics. *SIAM Rev.*, 51:3–33, 2009.

[13] A. Greenleaf, Y. Kurylev, M. Lassas, and G. Uhlmann. Approximate quantum and acoustic cloaking. *Journal Spectral Theory*, pages 27–80, 2011.

[14] A. Greenleaf, M. Lassas, and G. Uhlmann. Anisotropic conductivities that cannot be detected by EIT. *Physiol. Meas.*, 24:413–419, 2003.

[15] A. Greenleaf, M. Lassas, and G. Uhlmann. On nonuniqueness for Calderon’s inverse problem. *Math. Res. Lett.*, 10:685–693, 2003.

[16] M. Ikehata. Identification of the curve of discontinuity of the determinant of the anisotropic conductivity. *J. Inverse and Ill-posed Problems*, 8:273–285, 2000.

[17] V. Isakov. On uniqueness of recovery of discontinuous conductivity coefficient. *Commun. Pure Appl. Math.*, 41:865–877, 1988.
[18] H. Kang and J. K. Seo. The layer potential technique for the inverse conductivity problem. *Commun. Partial Diff. Eq.*, 12:267–278, 1996.

[19] R. Kohn and M. Vogelius. Determining conductivity by boundary measurements. *Commun. Pure Appl. Math.*, 37:289–298, 1984.

[20] R. V. Kohn, D. Onofrei, M. S. Vogelius, and M. I. Weinstein. Cloaking via change of variables for the Helmholtz equation. *Comm. Pure Appl. Math.*, 63:973–1016, 2010.

[21] R. V. Kohn, H. Shen, M. S. Vogelius, and M. I. Weinstein. Cloaking via change of variables in electrical impedance tomography. *Inverse Problems*, 24:015016, 2008.

[22] K. Kwon. Identification of anisotropic anomalous region in inverse problems. *Inverse Problems*, 20:1117–1136, 2004.

[23] K. Kwon and D. Sheen. Anisotropic inverse conductivity and scattering problems. *Inverse Problems*, 18:745–56, 2002.

[24] Y. Lai, J. Ng, H. Y. Chen, D. Z. Han, J. J. Xiao, Z.-Q. Zhang, and C. T. Chan. Illusion optics: The optical transformation of an object into another object. *Physical review letters*, 102:253902, 2009.
[25] M. Lassas and G. Uhlmann. On determining a Riemannian manifold from the Dirichlet-to-Neumann map. *Ann. Sci. Éc. Norm. Sup.*, 34:771–787, 2001.

[26] J. Lee and G. Uhlmann. Determining anisotropic real-analytic conductivities by boundary measurements. *Commun. Pure Appl. Math.*, 42:1097–1112, 1989.

[27] U. Leonhardt. Optical conformal mapping. *Science*, 312:1777–1780, 2006.

[28] U. Leonhardt and T. Tyc. Broadband invisibility by non-Euclidean cloaking. *Science*, 323:110–111, 2009.

[29] A. I. Nachmann. Reconstruction from boundary measurements. *Ann. Math.*, 128:531–576, 1988.

[30] A. I. Nachmann. Global uniqueness for a two-dimensional inverse boundary problem. *Ann. Math.*, 142:71–96, 1995.

[31] H. M. Nguyen. Cloaking via change of variables for the Helmholtz equation in the whole space. *Comm. Pure Appl. Math.*, 63:1505–1524, 2010.
[32] J. B. Pendry, D. Schurig, and D. R. Smith. Controlling electromagnetic fields. *Science*, 312:1780–1782, 2006.

[33] Z. Sun and G. Uhlmann. Anisotropic inverse problems in two dimensions. *Inverse Problems*, 19:1001–1010, 2003.

[34] J. Sylvester. An anisotropic inverse boundary value problem. *Commun. Pure Appl. Math.*, 43:201–232, 1990.

[35] J. Sylvester and G. Uhlmann. A global uniqueness theorem for an inverse boundary value problem. *Ann. math.*, 125:153–69, 1987.

[36] J. Sylvester and G. Uhlmann. Inverse boundary problem at the boundary - continuous dependence. *Commun. Pure Appl. Math.*, 41:197–219, 1988.