AlPd$_{15}$B$_7$: a new superconducting cage-compound with an anti-Yb$_3$Rh$_4$Sn$_{13}$-type of structure†

Qiang Zheng,† Walter Schnelle,† Yuri Prots,† Matej Bobnar,† Ulrich Burkhardt,† Andreas Leithe-Jasper† and Roman Gumeniuk* †a,b

A new intermetallic compound AlPd$_{15}$B$_7$ was synthesized by arc-melting the stoichiometric mixture of the elements. Single crystal X-ray diffraction data of ternary metal-rich boride reveal a new type of structure with the space group $Ia3d$ and the lattice parameter $a = 16.4466(3)$ Å. It adopts a filled anti-Yb$_3$Rh$_4$Sn$_{13}$-type structure, where the positions corresponding to 3Yb, 4Rh and 13Sn atoms are occupied by 3Pd, 4B, and 1Al + 12 Pd, respectively and 3B additionally at interstitial sites. Magnetic susceptibility, electrical resistivity, and specific heat measurements reveal bulk superconductivity with a critical temperature $T_c \approx 2.9$ K. Electronic structure calculations show that Pd 4d and B 2p states dominate the density of states (DOS) at the Fermi level $E_F$.

The variety of M, T, and X elements forming such compounds without major changes in the crystal structures, substitution by elements not only limited to the above mentioned groups should be possible.

In this study, we report the boride AlPd$_{15}$B$_7$ with a new type of structure, which is closely related to the Yb$_3$Rh$_4$Sn$_{13}$ prototype. Physical property measurements reveal superconducting behavior of this new boride with $T_c \approx 2.9$ K. The thermodynamic properties are compatible with an $s$-wave energy gap conventional superconductivity and moderate electron–phonon coupling.

Experimental section

Sample preparation

The samples with a nominal composition of AlPd$_{15}$B$_7$ were synthesized by arc-melting of aluminum (Chemapur, 99,9999 mass%) and palladium (Chemapur, 99.9 mass%) foils and crystalline boron powder (Alfa Aesar, 99.999 mass%). The obtained button was remelted several times to obtain a homogeneous sample. The weight loss during arc-melting was less than 0.1 mass%. All above operations were performed inside an argon-filled glove box ($p$(O$_2$/H$_2$O) ≤ 1 ppm). A nearly single phase sample was obtained. The resulting sample is stable in air for long time.

Powder and single-crystal X-ray diffraction

Due to the ductility, stress annealing of ground powders was performed at 840 °C for 1.5 h to obtain powder diffraction patterns with sharper and better resolved diffraction peaks. Powder X-ray diffraction (XRD) data were collected on a
HUBER G670 imaging plate Guinier camera (Cu $K_{α1}$ radiation, $\lambda = 1.540598$ Å). Phase analysis and indexing were performed within the WinXpow program package. Lattice parameters were refined by least-squares fitting with LaB$_6$ internal standard correction within the program package WinCSD.16 Single crystals were selected from the stress-annealed crushed sample. Single crystal XRD data were collected on a Rigaku AFC7 diffraction system equipped with a Saturn 724+ CCD detector (Mo $K_{α2}$ radiation, $\lambda = 0.71073$ Å). Absorption correction was made using a multi-scan procedure. The crystal structure was solved by a direct phase determination method and refined by a full-matrix least-squares procedure within the program package WinCSD. Details of the single crystal XRD data collection are listed in Table 1.

### Crystal structure determination

As shown in Fig. 1 and Fig. S1 (in the ESI†), the XRD patterns of the ground powders of an as-cast AlPd$_{15}$B$_7$ sample with and without stress-annealing reveal no difference between them besides peaks broadening in the latter one. As revealed by Fig. S2 in the ESI† the content of impurities in the as-cast sample is less than 1 vol%. The strong reflections could be indexed in a simple cubic unit cell with lattice parameter $a = 8.2233(2)$ Å. However, as one can see as an example in the inset of Fig. 1, there are still many reflections that could not be indexed in this unit cell, revealing the appearance of a superstructure. Taking these reflections into account, a body centered cubic unit cell with lattice parameter $a = 16.4466(3)$ Å is obtained. Subsequently, single crystals were selected for XRD experiments.

Details on single crystal XRD data collection for AlPd$_{15}$B$_7$ are listed in Table 1. Analysis of the reflection intensities and the extinction conditions confirm the appearance of the superstructure, indicating centrosymmetry and only one possible space group $Ia3d$. Positions of the heavy Pd and Al atoms and the light B atoms were acquired by a direct phase determination procedure and from difference Fourier maps, respectively.

### Results and discussion

#### Crystal structure determination

The magnetization data of a polycrystalline AlPd$_{15}$B$_7$ sample were collected in external fields $\mu_0H$ ranging from 2 mT to 7 T and temperatures between 1.8 K and 400 K on a SQUID magnetometer (MPMS-XL7, Quantum Design). The electrical resistivity was measured using a four-point ac method between 1.8 K and 320 K (PPMS, Quantum Design) on a small bar (1 mm × 1 mm × 6.5 mm). The heat capacity was measured by means of a relaxation-type calorimeter on a PPMS in external fields $\mu_0H$ up to 0.5 T between 0.35 and 320 K.

#### Electronic structure calculations

The scalar-relativistic band-structure of AlPd$_{15}$B$_7$ was calculated within the local-density approximation (LDA) of density functional theory (DFT) using the full-potential local-orbital FPLO code (version 9.01-35). The calculation was performed by employing the exchange-correlation potential of Perdew and Wang and using the experimental structural data from Table 2. The first Brillouin zone was sampled by a mesh of $12 \times 12 \times 12$ (1728) $k$ points.

### Table 1. Crystallographic data for AlPd$_{15}$B$_7$

| Composition | Space group | Ia3d |
|-------------|-------------|------|
| Lattice parameters | $a$ (Å) | 16.4466(3) |
| | $V$ (Å$^3$) | 4448.7(3) |
| Formulation per cell, Z | 16 |
| Calculated density (g cm$^{-3}$) | 10.14 |

### Table 2. Atomic coordinates and isotropic displacement parameters (in Å$^2$) for AlPd$_{15}$B$_7$

| Atom | Site | $x$ | $y$ | $z$ | $B_{iso}$/Eq$^a$ |
|------|------|-----|-----|-----|-----------------|
| Al   | 16a  | 0   | 0   | 0   | 1.08(8)         |
| Pd1  | 48g  | 1/8 | 0.24484(6) | 0.00516(6) | 1.23(2)         |
| Pd2  | 96h  | 0.00099(6) | 0.14693(6) | 0.07718(6) | 0.99(2)         |
| Pd3  | 96h  | 0.10758(6) | 0.24514(6) | 0.33770(6) | 1.13(2)         |
| B1   | 16b  | 1/8 | 1/8 | 1/8 | 1.4(5)          |
| B2   | 48f  | 0.263(2) | 0   | 1/4 | 1.5(3)          |
| B3   | 48g  | 1/8 | 0.619(1) | 0.631(1) | 1.4(3)          |

$^a$ $B_{iso} = 1/3[a^2+2b^2+2c^2] + 2abc(\cos\gamma)B_{123} + 2aca(\cos\beta) B_{13} + 2bcba(\cos\alpha)B_{12}$.
Anisotropic displacement parameters for all the Al and Pd atoms were refined. Final atomic coordinates and anisotropic atomic displacement parameters are listed in Tables 2 and 3, respectively. Moreover, diffraction data were also collected for single crystals selected from powders of the as-cast sample without stress annealing, and reveal the same structure solution result and only larger residual values due to broadening of the reflections. The acquired composition is in good agreement with the compositions from the EDXS and WDXS analysis.

Crystal structure

AlPd15B7 crystallizes with a new type of structure, which can be regarded as a filled anti-Yb3Rh4Sn13 structure. The crystal structures of AlPd15B7 and Yb3Rh4Sn13 are shown in Fig. 2a and Fig. S3 in the ESI† respectively. The group–subgroup relationship between them is shown in Fig. 2b. The crystallographic positions for 3Yb, 4Rh and 13Sn atoms in the Yb3Rh4Sn13 structure1 are now occupied by 3Pd, 4B, and 1Al + 12Pd atoms in the AlPd15B7 structure, respectively. The octahedra formed by Yb and Sn atoms are empty in the Yb3Rh4Sn13-type structure, whereas the corresponding octahedra solely formed by Pd atoms are filled by additional B atoms in the AlPd15B7 structure. The transformation between these two types of structures can be described as:

\[
M_3T_4X_{12}X_1 \rightarrow M_1T_15X_7
\]

where M, T, and X are Yb, Rh, and Sn in Yb3Rh4Sn13, and Al, Pd and B in AlPd15B7. Different from Yb3Rh4Sn13, the lattice parameters of AlPd15B7 are doubled and the space group changes from \( \text{Pm} \overline{3n} \) to \( \text{Ia} \overline{3d} \), due to the shift of the Pd atoms from some special positions. As one can see in Fig. 6 in ref. 13, the relationship between many derivatives of the primitive cubic Yb3Rh4Sn13-type of structure (e.g. \( \text{La}_3\text{Rh}_4\text{Sn}_{13} \) (I\( \text{4} \)\( \text{3} \)), \( \text{Ca}_3\text{Pt}_4\text{Ge}_{13-y} \) (I\( \text{2} \)\( \text{3} \))),12 \( \text{HT-Y}_3\text{Pt}_4\text{Ge}_{13} \) (R3\( \text{c} \)),13 \( \text{Yb}_3\text{Pt}_4\text{Ge}_{13} \) (P4\( \text{2} \)cm),12 and \( \text{LT-Y}_3\text{Pt}_4\text{Ge}_{13} \) (C\( \text{c} \))14) can be understood within a group–subgroup scheme. AlPd15B7 crystallizes with the space group \( \text{Ia}3d \), which links the space groups \( \text{Pm}3n \) and \( \text{Ia}3 \) (Fig. 3).

### Table 3

| Atom | \( B_{11} \) | \( B_{12} \) | \( B_{13} \) | \( B_{12} \) | \( B_{13} \) | \( B_{13} \) |
|------|----------|----------|----------|----------|----------|----------|
| Al   | 1.1(1)   | 1.11(1)  | -0.06(1) | 0.02(2)  | 0.12(4)  | 0.12(4)  |
| Pd1  | 1.09(5)  | 1.31(5)  | -0.11(3) | 0.12(4)  | 0.01(2)  | 0.02(2)  |
| Pd2  | 0.93(3)  | 1.09(3)  | 0.96(4)  | -0.09(3) | -0.02(3) | -0.14(3) |
| Pd3  | 0.26(4)  | 1.06(4)  | 1.10(4)  | -0.06(3) | -0.02(3) | -0.14(3) |

Fig. 1 Experimental (black circles) and calculated (red line) powder XRD patterns of an as-cast AlPd15B7 sample after grinding and subsequent stress-annealing at 840 °C for 1.5 h. Peak positions of AlPd15B7 are given by black ticks; the difference plot is shown as a black line in the bottom. The inset shows the 2\( \theta \)-range from 45.8° to 52°, where the reflections corresponding to the superstructure are marked by asterisks (see text).

Fig. 2 (a) Crystal structure of AlPd15B7. For a better visualization, only a part of the structure in one unit cell is shown. [Al@Pd12] icosahedra – light gray; distorted [B@Pd6] trigonal prisms – dark gray; [B@Pd6] octahedra – blue. (b) group–subgroup relationship between the structures of Yb3Rh4Sn13 and AlPd15B7 (see text). For better visualization the coordinates for AlPd15B7 are not standardized, like those given in Table 2.
the indices of $k4$ (klassengleiche) and $t2$ (translationengleiche) symmetry reductions, respectively.

The AlPd$_{15}$B$_7$ structure is shown in Fig. 2a, and interatomic distances are listed in Table S1 in the ESI. In such a structural arrangement, similar to Al residing in [Ni$_{12}$] icosahedra in the AlNi$_{6}$B$_{8}$ structure, each Al atom resides in a distorted [Pd$_{12}$] icosahedron, consisting of 6Pd2 and 6Pd3 atoms, with the Al–Pd and Pd–Pd distances ranging from 2.730 Å to 2.752 Å, and from 2.714 Å to 2.955 Å, respectively. The empty spaces between the icosahedra are occupied by the Pd$_1$, B$_1$, B$_2$ and B$_3$ atoms. Each [Al@Pd$_{12}$] icosahedron interconnects with 14 neighboring [Al@Pd$_{12}$] icosahedra via 8 distorted corner-sharing trigonal prisms and 6 distorted octahedra. The two sorts of corner-sharing [Pd$_n$] trigonal prisms, formed by 6 Pd2 atoms or 2Pd2 and 4Pd3 atoms, are centered by the B1 or B3 atoms, respectively. In a [B1Pd$_6$] trigonal prism, the Pd2–Pd2 distances are 2.850 Å and 2.942 Å, and the Pd3–B3 distances are 2.166 Å to 2.180 Å; in Pd2B (Fe$_3$C type of structure), B resides in trigonal prisms with Pd–B distances of 2.166 Å to 2.180 Å; in a [B3Pd$_6$] trigonal prism, the Pd–Pd distances are in the range of 2.885 Å to 2.955 Å, the Pd2–B3 distance is 2.139 Å, and the Pd3–B3 distances are 2.157 Å and 2.28 Å. Each [Pd$_8$] octahedron, formed by Pd2, Pd2, and 2Pd3 atoms, is centered by a B2 atom, with Pd–Pd distances and Pd–B distances ranging from 2.714 Å to 3.5423 Å and from 2.08 Å to 2.252 Å, respectively. These kinds of metal–atom coordination of boron are commonly observed in borides. In Pd$_2$B$_7$ (Fe$_3$C type of structure), B resides in trigonal prisms with Pd–B distances of 2.166 Å to 2.180 Å; in [Pd$_3$B$_2$] (anti-CaCl$_2$ type of structure), B centers Pd octahedra with Pd–B distances of 2.104 Å to 2.122 Å. In comparison with the sum of atomic radii of pure elements ($r_{\text{Al}} = 0.83 \, \text{Å}$, $r_{\text{Pd}} = 1.43 \, \text{Å}$, and $r_{\text{Pd}} = 1.38 \, \text{Å}$ [ref. 24]), the Al–Pd distances are shorter than the sum of atomic radii of Al and Pd, the Pd–Pd distances obviously deviate from the two-fold atomic radius of Pd, and the shortest and longest Pd–B distances are 6% shorter and ~3% longer than the sum of atomic radii of Pd and B, respectively. Coordination polyhedra of all the atoms in AlPd$_{15}$B$_7$ are shown in Fig. S4 in the ESI.

**Physical properties**

The magnetic susceptibility of AlPd$_{15}$B$_7$ in an external field of $\mu_0 H = 7$ T is shown in the inset of Fig. 3, revealing that it is a diamagnet with an extrapolated value of $\chi_0 = -2.0 \times 10^{-4}$ emu mol$^{-1}$ at $T = 0$. The slight upturn towards low temperature is probably due to minor paramagnetic impurities or point defects.

The temperature dependence of the zero-field-cooled (zfc) and field-cooled (fc) magnetic susceptibility of AlPd$_{15}$B$_7$ in an external field of $\mu_0 H = 2$ mT is shown in Fig. 3, indicating a superconducting transition. The transition temperature was determined via a tangent to the steepest slope of $\chi_{\text{zfc}}(T)$, resulting in $T_{c,\text{zfc}} = 2.78$ K. Considering the demagnetization correction, the diamagnetic response in zfc is close to complete. The fc signal (Meißner effect) is much weaker, which is due to strong flux line pinning in the type-II superconductor (see below).

As shown in Fig. 4, the electrical resistivity of AlPd$_{15}$B$_7$ increases with increasing temperature in the range of 3.0–320 K, revealing a metallic behavior of this compound. For $4 \, K < T < 40 \, K$, it can be fitted to $\rho_0 + A_{\text{FL}} T^2$ with a residual resistivity $\rho_0 = 0.20(1) \, \mu\Omega \cdot \text{m}$ from impurity scattering, and the cross section of quasiparticle–quasiparticle scattering $A_{\text{FL}} = 3.96(3) \times 10^{-5} \, \mu\Omega \cdot \text{m} \cdot \text{K}^{-2}$. The room-temperature resistivity $\rho(300 \, K)$ is 0.68 $\mu\Omega \cdot \text{m}$, and the residual resistivity ratio $\text{RRR} = 3.4$ indicates moderate quality of the investigated polycrystalline sample. As revealed by the inset of Fig. 4, a sharp superconducting transition can be observed, and the resistivity drops to zero at $T_c = 2.89$ K.

The specific heat $c_p/T$ vs. $T$ for AlPd$_{15}$B$_7$ in the temperature range 0.35–4.0 K and in various magnetic fields is shown in Fig. 4.
Fig. 5a. For zero-field, a sizable sharp step-like anomaly confirms the bulk nature of the superconductivity. The normal-state specific heat (Fig. 5b) can be described by $c_{p,n}(T) = \gamma_{tot}T + \beta T^2 + \delta T^3$ in the temperature range 0.5–5 K, where $\gamma_{tot}$ is the Sommerfeld coefficient of the electronic heat capacity, and $\beta T^2 + \delta T^3$ are the first terms of the harmonic lattice approximation for the phonon contribution. The resulting parameters are $\gamma_{tot} = 28.3(1) \text{ mJ mol}^{-1} \text{ K}^{-2}$ and $\beta = 1.38(2) \text{ mJ mol}^{-1} \text{ K}^{-3}$, which corresponds to an initial Debye temperature $\Theta_D(0) = 319$ K, and $\delta = 2.86(6) \times 10^{-5} \text{ mJ mol}^{-1} \text{ K}^{-6}$.

The electronic contribution $c_{el}(T, H)$ below $T_c$ can be analyzed by subtracting the lattice terms $\beta T^2 + \delta T^3$. Fig. 5c shows $\Delta c_p = c_{el} - \gamma_{tot}T$ which is well described by the BCS theory (for weak electron–phonon coupling) in the whole temperature range. The resulting parameters are the jump $\Delta c_p/T_c = 37.5 \text{ mJ mol}^{-1} \text{ K}^{-2}$, the transition midpoint $T^*_c = 2.796$ K, and the width of the transition of 0.026 $T_c$. The ratio $\Delta c_p/(\gamma_{tot}T_c)$ of 1.32 is somewhat lower than the value 1.43 from the BCS theory. At the lowest temperatures (in the range 0.35–0.70 K) $c_{el}(T)$ is well fitted by $c_{el} = \gamma_0 T + \gamma T^* T^n e^{-s_T^* T}$ (as shown in Fig S5a in the ESI†). Such exponential behavior is expected for an s-wave superconductor without nodes of the gap, as in the BCS theory. The resulting parameters for zero field are $\gamma_0 = 2.17 \text{ mJ mol}^{-1} \text{ K}^{-2}$ and the energy gap ratio $\Delta/k_B T_c = 1.08$. The non-zero $\gamma_0$ is probably due to the presence of a non-superconducting metallic impurity phase. We may thus correct the normal state electronic specific heat coefficient by $\gamma_n = \gamma_{tot} - \gamma_0$. The corrected value for the jump $\Delta c_p/(\gamma_n T_c)$ = 1.14 reproduces even better the weak-coupling BCS value 1.43. However, the energy gap ratio $\Delta/k_B T_c$ observed at the lowest temperatures is much smaller than expected from BCS theory ($\Delta/k_B T_c = 1.76$). It is therefore possible that AlPd$_{13}$B$_7$ needs to be described within a two-gap model. The resolution of the present specific heat data, due to the small available mass for the study, is insufficient for a more sophisticated analysis.

The DFT calculations reveal a DOS at the Fermi level $N(E_F) = 7.9$ states eV$^{-1}$ f.u.$^{-1}$ (see below), converting to a “bare” Sommerfeld coefficient of the electronic specific heat $\gamma_{bare} = 18.6 \text{ mJ mol}^{-1} \text{ K}^{-2}$. The electron–phonon coupling parameter $\lambda$ in the relation $\gamma_n = (1 + \lambda) \gamma_{bare}$ is calculated to be $0.40$. $\lambda$ can also be estimated from the McMillan formula $\ln(1.45 T_c/\Theta_D) = -1.04 (1 + \lambda) \ln(1 - \mu^* (1 + 0.62\lambda))$. Assuming the common value $\mu^* = 0.15$, $\lambda$ results as 0.58, which is somewhat larger than the above value.

The fitted Sommerfeld parameters $\gamma(H) = \gamma_0$, extrapolated to $T = 0$, were plotted against the fields (Fig. S5b in the ESI†), showing a linear relationship with the magnetic fields. This also indicates a typical s-wave gap for AlPd$_{13}$B$_7$. The upper critical field values were estimated from the midpoints of superconducting transition in $c_p(T, \mu_0 H)$, as shown in Fig. 5d. It is obvious that $T_c(\mu_0 H)$ varies almost linearly with $\mu_0 H$. This most extreme extrapolation results in $\mu_0 H_c(0) \approx 500(100)$ mT.

**Electronic structure**

The total and partial (atomic and orbital resolved) density of states (DOS) for AlPd$_{13}$B$_7$ are shown in Fig. 6(a–c). Similar to the electronic structures of the M$_3$Pd$_3$Ge$_{11}$ compounds,$^{12,13}$ the Fermi level $E_F$ for AlPd$_{13}$B$_7$ is located close to a dip in the DOS. Thus, the system is a metal with 7.9 states eV$^{-1}$ f.u.$^{-1}$ at $E_F$. 

![Fig. 5](image-url) (a) Molar specific heat $c_p/T$ of AlPd$_{13}$B$_7$ in various magnetic fields; (b) the low-temperature normal-state $c_p(T)$ for AlPd$_{13}$B$_7$ together with the fit to $c_p = \gamma_{tot}T + \beta T^2 + \delta T^3$; (c) fitting of the difference specific heat $\Delta c_p = c_{el} - \gamma_{tot}T$ based upon the BCS theory; (d) upper critical field $\mu_0 H_c$ of AlPd$_{13}$B$_7$ vs. $T$ derived from the midpoints of the superconducting transitions in $c_p(T, H)$.

![Fig. 6](image-url) (a) Total and atom resolved electronic density of states (DOS) for AlPd$_{13}$B$_7$; (b) orbital resolved DOS for (b) palladium and (c) boron atoms.
This DOS is mostly due to the contributions of the Pd 4d (~71%) and the B 2p (~25%) states. The broad valence band (~7 eV) in the electronic structure of AlPd15B7 is dominated by Pd 4d states (Fig. 6b) with sizeable B 2p contributions (Fig. 6c) at lower energies (between ~7 eV and ~3.5 eV). Furthermore, the lowest lying separate band (between ~11 eV and ~7.5 eV) is mainly formed by B 2s and Pd 4d states. All this is related to the covalent bonding in the Pd–B framework. Al contributions to the valence band are negligibly small. Interestingly, a narrow separate band (~0.3 eV) centered at ~7.1 eV is mostly due to 4d states of Pd2 and Pd3 as well as Al 3s. Obviously, it reflects covalent interaction of Al and Pd atoms within the [Al@Pd12] icosahedra (Fig. 2). The electronic band structure of AlPd15B7 is reminiscent of the electronic structures of such Pd-rich borides as PdB and SrPdB. They also possess a broad valence band dominated by Pd 4d and B 2p states followed by a separate low energy band due to the mixing of B 2s and Pd 4d states. The differences are that for AlPd15B7 a larger DOS at EF as well as an additional narrow band at ~7.1 eV are observed.

Conclusions

We gratefully acknowledge Prof. Yu. Grin’s interest and steady support. A new ternary boride AlPd15B7, with a filled anti-Yb3Rh4Sn13-type structure, has been synthesized by arc-melting mixtures of elements, and its crystal structure was determined from the single-crystal X-ray data. The crystal structure of AlPd15B7 at a heat measurements show a superconducting transition of Tc of about 2.9 K. The specific heat measurements reveal AlPd15B7 to be a conventional s-wave superconductor with weak electron–phonon coupling according to the BCS theory. Electronic structure calculations by DFT reveal that the DOS at EF is dominated by d states.

AlPd15B7 is the first boride with a crystal structure comprising the framework of the Remeika Yb3Rh4Sn13 prototype but being decorated by completely different types of elements. The occurrence of this new type of structure suggests that the constituting elements in the Remeika structure framework could be extended to a variety of elements in future investigations.

Acknowledgements

We thank Mr S. Hückmann and Dr H. Borrmann for performing powder X-ray diffraction measurements, Ms. P. Scheppan, Ms. M. Eckert and Ms. S. Kostmann for the metallographic analysis, and Mr R. Koban for physical property measurements.

References

1.  J. L. Hodeau, J. Chenavas, M. Marezio and J. P. Remeika, Solid State Commun., 1980, 36, 839.
2.  P. Villars, Pearson’s Handbook, Crystallographic Data for Intermetallic Phases, Materials Park, OH, 1997.
3.  J. M. Vandenbergh, Mater. Res. Bull., 1980, 15, 835.
4.  S. Miraglia, J. L. Hodeau, M. Marezio, C. Laviron, M. Ghedira and G. P. Espinosa, J. Solid State Chem., 1986, 63, 358.
5.  G. K. Shenoy, B. D. Dunlap and F. Y. Fradin, Ternary Superconductors, Elsevier, North-Holland, Amsterdam, 1981.
6.  J. P. Remeika, G. P. Espinosa, A. S. Cooper, H. Barz, J. M. Rowell, D. B. McWhan, J. M. Vandenbergh, D. E. Moncton, Z. Fisk, L. D. Woolf, H. C. Hamaker, M. B. Maple, G. Shirane and W. Thomson, Solid State Commun., 1980, 34, 923.
7.  Y. Mudryk, A. Grytsiv, P. Rogl, C. Dusek, A. Galatanu, E. Idl, H. Michor, E. Bauer, C. Godart, D. Kaczorowski, L. Romaka and O. Bodak, J. Phys.: Condens. Matter, 2001, 13, 7391.
8.  N. Kase, H. Hayamizu and J. Akimitsu, Phys. Rev. B: Condens. Matter, 2011, 83, 184509.
9.  B. K. Rai, I. W. H. Oswald, J. K. Wang, G. T. McCandless, J. Y. Chan and E. Morosan, Chem. Mater., 2015, 27, 2488.
10.  W. Jeitschko and D. Braun, Acta Crystallogr., 1977, 33, 3401.
11.  R. Gumeniuk, H. Borrmann, A. Ormecki, H. Rosner, W. Schnelle, M. Nicklas, Yu. Grin and A. Leithe-Jasper, Z. Kristallogr., 2010, 225, 531.
12.  R. Gumeniuk, L. Akselrud, K. O. Kvashnina, W. Schnelle, A. A. Tsirlin, C. Curfs, H. Rosner, M. Schöneich, U. Burkhardt, U. Schwarz, Yu. Grin and A. Leithe-Jasper, Dalton Trans., 2012, 41, 6299.
13.  R. Gumeniuk, M. Schöneich, K. O. Kvashnina, L. Akselrud, A. A. Tsirlin, M. Nicklas, W. Schnelle, O. Janson, Q. Zheng, C. Curfs, U. Burkhardt, U. Schwarz and A. Leithe-Jasper, Dalton Trans., 2015, 44, 5638.
14.  R. Gumeniuk, M. Nicklas, L. Akselrud, W. Schnelle, U. Schwarz, A. A. Tsirlin, A. Leithe-Jasper and Yu. Grin, Phys. Rev. B, 2013, 87, 224502.
15.  WinXpow, STOE and Cie GmbH, Darmstadt, 2003.
16.  L. Akselrud and Yu. Grin, J. Appl. Crystallogr., 2014, 47, 803.
17.  K. Koepernik and H. Eschrig, Phys. Rev. B, 1999, 59, 1743.
18.  J. P. Perdew and Y. Wang, Phys. Rev. B, 1992, 45, 13244.
19.  Q. Zheng, F. R. Wagner, A. Ormecki, Yu. Prots, U. Burkhardt, M. Schmidt, W. Schnelle, Yu. Grin and A. Leithe-Jasper, Chem. – Eur. J., 2015, 21, 16532.
20.  Yu. B. Kuz’ma, Crystallochemistry of Borides, Livv University Publishers, Livv, 1983.
21.  P. Rogl, in Inorganic Reactions and Methods, ed. A. P. Hagen, Wiley, 1991.
22.  E. Stenberg, Acta Chem. Scand., 1961, 15, 861.
23.  L.-E. Tergenius and T. Lundström, J. Solid State Chem., 1980, 31, 361.
24.  J. Emsley, The Elements, Clarendon Press, Oxford, 1998.
25.  R. Gumeniuk, M. Schmitt, W. Schnelle, U. Burkhardt, H. Rosner and A. Leithe-Jasper, Z. Anorg. Allg. Chem., 2010, 636, 954.