Extended Molecular Dynamics Methods for Vortex Dynamics in Nano-structured Superconductors

Masaru Kato\textsuperscript{a,}\textsuperscript{*}, Osamu Sato\textsuperscript{b}

\textsuperscript{a}Department of Mathematical Sciences, Osaka Prefecture University, 1-1, Gakuencho, Naka-ku, Sakai, Osaka 599-8531, Japan
\textsuperscript{b}Osaka Prefecture University College of Technology, 26-12, Saiwaicho, Neyagawa, Osaka 572-8572, Japan

Abstract

Using improved molecular dynamics simulation method, we study vortex dynamics in nano-scaled superconductors. Heat generations during vortex motion, heat transfer in superconductors, and entropy forces to vortices are incorporated. Also quasiparticle relaxations after vortex motion, and their attractive “retarded” forces to other vortices are incorporated using the condensation-energy field. We show the time development of formation of vortex channel flow in a superconducting Corbino-disk.

1. Introduction

For applications of superconductors, vortex dynamics is a key feature, because when vortices move, the zero resistivity of a superconductor is broken. In order to investigate the dynamics of vortices theoretically or by computer simulations, we use phenomenological time-dependent Ginzburg-Landau equations [1] or the molecular dynamics method. Especially, the molecular dynamics method is effective when we consider many vortices in superconductors [2], because we only consider classical equations of motion of the vortices. However, in this method, several features are missing, when we consider real superconducting systems. For example, when vortices move, heat generation occurs. If the vortex motion is uniform, this heat generation is also uniform. But when vortex motion is not uniform, such as in a corbino disk [3], non-uniform heat generation occurs and then non-uniform temperature distribution appears. A vortex has a transport-entropy [4,5], and then motion of the vortex is affected by this temperature distribution. In order to incorporate such effect, we must solve the heat transport equation with the molecular dynamics equation for vortices. Another feature is a retardation effect, which comes from quasiparticle recombination after fast movement of vortices [6,7]. After a vortex moves, the order parameter that was inside of the vortex core is restored to the uniform value, but it takes a time for recombining Cooper pairs. Then, if vortices move fast, the vortex motion is affected by a preceding vortex. In order to incorporate such retarded effect, we introduce a condensation energy field, which is proportional to the square of the absolute value of the order parameter. In Sec. 2, we explain our method, and in Sec. 3 we show our numerical results.

* Corresponding author. Tel.: +81-72-254-9368; fax: +81-72-254-9916.
E-mail address: kato@ms.osakafu-u.ac.jp

© 2015 The Authors. Published by Elsevier B.V. This is an open access article under the CC BY-NC-ND license (http://creativecommons.org/licenses/by-nc-nd/4.0/).
Peer-review under responsibility of the ISS 2014 Program Committee

Keywords: Vortex dynamics; molecular dynamics; heat transport; quasiparticle realxation; Corbino-disk; finite element method
2. Method

In usual molecular dynamics method, we ignore the inertia mass term because of dissipative dynamics of vortices and then we treat the overdamped equations of motions, where a damping term is large. We denote the position of i-th vortex \( r_i \). Forces on this vortex are, the vortex-vortex interaction force \( f_{vv}^{\text{mv}} \), and the force of vortex interaction with the Meissner shielding current \( f_{H}^{\text{mv}} \), the pinning force from the impurities \( f_{\text{pin}}^{\text{imp}} \), and a driving force from an external current \( f_d \). Also thermal fluctuation force \( f_{\text{fl}} \) from the environment is incorporated and the equation of motion is based on the Langevin dynamics as follows,

\[
\eta \frac{dn}{dt} = f_d + f_{vv}^{\text{mv}} + f_{H}^{\text{mv}} + f_{\text{pin}}^{\text{imp}} + f_{\text{fl}},
\]

where \( \eta \) is the viscosity. The driving, pinning and fluctuation forces are defined as,

\[
f_d = \frac{1}{c} j(r_i) \times \Phi_0,
\]

\[
f_{vv}^{\text{mv}} = \sum_j \frac{f_{ij}}{r_{ij}} \Theta \left( \frac{r_i - r_j}{\lambda} \right) \delta_{ij},
\]

\[
\left\langle f_i'(t_1), f_j'(t_2) \right\rangle = 2 \eta k_BT \delta(t_1-t_2),
\]

where \( j(r_i) \) is an external current at the vortex position \( r_i \), \( r_j \) is the position of j-th impurity, \( r_{ij} \) is the size of impurities, \( f_s \) is the strength of pinning force, \( \lambda \) is the penetration depth, and \( \Theta(x) \) is a step function. Thermal fluctuation force depends on the temperature \( T \). In Sec. 3, we show numerical simulation results for a superconducting corbino-disk with radius \( R \), and in such geometry, the vortex-vortex interaction force and the force from shielding current are given as,

\[
f_{vv}^{\text{mv}} = f_0 \sum_j \left( \frac{\hat{r}_{ij} \cdot r_i}{r_{ij}} - \frac{r_i^2}{R} \right) \left( \frac{r_i^2}{R^2} - r_j^2 \right),
\]

\[
f_{H}^{\text{mv}} = f_0 \left( \frac{1}{1 - (r_i/R)^2} \right) \frac{\pi R^2 H}{\Phi_0} r_i,\]

where \( f_0 \) is the strength of vortex-vortex interaction. In previous our work [8], we use a vortex-vortex interaction in 3-dimensional superconductors. But in this work, we use a vortex-vortex interaction in 2-dimensional superconductors.

For the heat generation by the vortex motion, we must solve the heat transport equation and we use the finite element method (FEM),

\[
\sum_i (a_i T_i + c_i T_i) = Q_j,
\]

where temperature is expanded as \( T(x,y,t) = \sum_i T_k(t) N_k(x,y) \), using area coordinates \( N_k(x,y) \). And \( c_g = \alpha_I \cdot c + \lambda h (K_{\rho}^{\text{mv}} + K_{\rho}^{\text{mm}}) \), \( a_g = c h l_0 \) and \( Q_j = \alpha T_a \sum_i G_k + \sum_i h a l_0 \). Here \( c \), \( \lambda \) and \( h \) are the specific heat, thermal conductivity, and he thickness of the superconductor, respectively. \( a \) and \( T_a \) are thermal conductance between the superconductor and substrate and temperature of the substrate, respectively. The heat generation by the vortex motion is given by,
\[ q(x,y) = \sum_j f_j(x_j,y_j) \cdot \frac{1}{S_e} \cdot v_j, \]  

(8)

where the summation is taken over vortices that are in the element of the FEM around the position \((x,y)\). From this temperature distribution, the entropy force on the vortex \(f_T = -S_g \text{grad}T\) occurs, because of entropy \(S_g\) of vortex.

When vortex moves fast, then it takes a time for recovering of the order parameter to the equilibrium value, or Cooper pair recombination after the vortex movement. In the GL equations, such effect can be treated [1], because the order parameter variation can be calculated explicitly. In order to incorporate this effect, we consider a condensation energy field, which is proportional to the GL free energy. A local equilibrium condensation energy at temperature \(T(r)\) is given as,

\[ E_{\text{cond}}(r) = E_{\text{cond}}^0 \left(1 - \frac{T(r)}{T_c}\right)^2, \]  

(9)

where \(E_{\text{cond}}\) is the condensation energy at \(T = 0\). After the vortex leaves, the condensation energy \(V_{\text{cond}}(r)\) decreases and recovers to the equilibrium value with the relaxation time \(\tau\) as,

\[ \frac{\partial V_{\text{cond}}(r)}{\partial t} = \frac{E_{\text{cond}}(r) - V_{\text{cond}}(r)}{\tau}. \]  

(10)

The vortex gets the force from this condensation energy field as \(f_{R_i} = -\text{grad}V_{\text{cond}}\). We take these forces in to the molecular dynamics equation.

3. Numerical Results

We consider a superconducting Corbino disk with \(R = 10\xi_0\), where an external magnetic field is applied to perpendicular to the disk and an external current flows from the center of the disk to the edge. Here \(\xi_0\) is the coherence length at \(T = 0\).

\[ \text{Fig. 1. Time development of condensation energy distribution and vortex configuration for } t/(\eta_0^2/\xi_0) = 0 \text{ (a), 0.01 (b), 1.0 (c), 2.0 (d), 8.0 (e), 9.8 (f).} \]

In Fig. 1, we show the time development of the vortices and the condensation energy distribution. Initial positions of vortices were given randomly. It is seen that initially vortices moves independently and the condensation energy is low.
around vortices (Figs. 1 (a) and (b)). But gradually vortices follow other vortices (Figs. (c), (d) and (e), and finally they make weak condensation energy circles (Fig. 1 (f)). These circles may be considered as flux flow channels.

In Fig. 2, the time development of the temperature distribution is shown. Initially, the temperature distribution is uniform (Fig. 2 (a)). Then around moving vortices, temperature becomes high (Fig. 2 (b)). Because vortex motion is fast around the center, the temperature becomes high around the center (Figs. 2 (c), (d) and (e)). There are also high temperature circles (Fig. 2 (f)), which coincide with weak condensation energy circles in Fig. 1 (f).

Fig. 2. Time development of temperature distribution and vortex configuration for $1/(\eta u/f_0) = 0$ (a), 0.01 (b), 1.0 (c), 2.0 (d), 8.0 (e), 9.8 (f).

4. Summary

We have developed the molecular dynamics method for vortex motion in superconductors. Using this method, we have investigated fast vortex dynamics in a superconducting corbino disk. We found vortices make flow channels where the condensation energy is low and the temperature is high.

Our method is applicable to other geometries, such as strip superconductors or square superconducting plates. We will investigate the vortex dynamics in these superconductors in future.

Acknowledgements

We thank T. Koyama, T. Ishida, T. Tamegai, S. Okuma, N. Kokubo, K. Hirata, M. Hayashi, H. Ebisawa for useful discussions. This work was supported by JSPS KAKENHI Grant Number 26400367.

References

[1] R. Kato, Y. Enomoto and S. Maekawa, Phys. Rev. B 47 (1993) 8016–8024.
[2] C. Reichhardt, C. J. Olson and F. Nori, Phys. Rev. Lett. 78 (1997) 2648-2651.
[3] S. Okuma, S. Morishima and M. Kamada, Phys. Rev. B 76 (2007) 224521(1-6).
[4] K. Maki, Prog. Theor. Phys. 41 (1969) 902-918.
[5] K. Maki, Phys. Rev. B 43 (1991) 1252-1254.
[6] D. Y. Vodolazov and F. M. Peeters, Phys. Rev. B 76 (2007) 014521(1-9).
[7] S. Okuma, D. Shimamoto and N. Kokubo, Phys. Rev. B 85 (2012) 064508 (1-5).
[8] D.E. Fujibayashi, M. Kato, Physica C 484 (2013) 94.