Measuring rare and exclusive Higgs boson decays into light resonances

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Abstract We evaluate the LHC’s potential of observing Higgs boson decays into light elementary or composite resonances through their hadronic decay channels. We focus on the Higgs boson production processes with the largest cross sections, $pp \rightarrow h$ and $pp \rightarrow h + \text{jet}$, with subsequent decays $h \rightarrow ZA$ or $h \rightarrow Z \eta_c$, and comment on the production process $pp \rightarrow hZ$. By exploiting track-based jet substructure observables and extrapolating to 3000 fb\textsuperscript{-1} we find $\mathcal{BR}(h \rightarrow ZA) \simeq \mathcal{BR}(h \rightarrow Z \eta_c) \lesssim 0.02$ at 95 % CL. We interpret this limit in terms of the 2HDM Type 1. We find that searches for $h \rightarrow ZA$ are complementary to existing measurements and can constrain large parts of the currently allowed parameter space.

1 Introduction

The greatly successful Run 1 of the large hadron collider (LHC) culminated in the discovery of a state that resembles the standard model (SM) Higgs boson [1,2]. First measurements of its couplings to gauge bosons and second-generation quarks [7,8]. However, for SM couplings the branching ratios for exclusive Higgs boson decays are generally of $\mathcal{O}(10^{-5})$ or less [7,9,10], e.g. $\mathcal{BR}(h \rightarrow Z \eta_c) \lesssim 1.4 \times 10^{-5}$, $\mathcal{BR}(h \rightarrow \rho^0 \gamma) \simeq 1.68 \times 10^{-5}$ or $\mathcal{BR}(h \rightarrow J/\psi \gamma) \simeq 2.95 \times 10^{-5}$, resulting in small expected event yields. Nevertheless, both general purpose experiments at the LHC have performed searches for exclusive Higgs boson decays, focusing on the dimuon decays of vector quarkonia. With Run 1 data the ATLAS collaboration obtained a similar upper limit for $\mathcal{BR}(h \rightarrow J/\psi \gamma)$ [11], while the CMS collaboration obtained a similar upper limit for $\mathcal{BR}(h \rightarrow J/\psi \gamma)$ [12]. Recently, the ATLAS collaboration has also set a 95 % CL upper limit of $1.4 \times 10^{-3}$ on $\mathcal{BR}(h \rightarrow \phi \gamma)$ [13].

Hence, rare decays of Higgs bosons into light elementary or composite resonances are of direct relevance for the two most important tasks of the upcoming LHC runs: (a) precision measurements of the Higgs boson properties; and (b) searches for new physics.

While most existing search strategies rely upon resonance decays into leptons, i.e. muons, the total width of most composite resonances and elementary scalars is dominated by decays into hadronic final states, e.g. $\mathcal{BR}(\eta_c \rightarrow...
hadrons) > 52%\textsuperscript{1} [14]. Instead of exploiting only leptonic decay modes, we therefore propose that the inclusive hadronic decays be considered. Light resonances \(X\) with masses of \(m_X = 1 - 10\) GeV produced in decays of the Higgs boson with a mass of 125 GeV, are highly boosted and their decay products are thus confined within a small area of the detector. The angular separation of the decay products of the resonance \(X\) scales like \(\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2} \sim 4m_X/m_H\), where \(\eta\) is the pseudorapidity and \(\phi\) the azimuthal angle. Separating the decay products in the calorimeters of the detector poses a challenge, as the typical size of hadronic calorimeter cells is \(0.1 \times 0.1\) in the \((\eta, \phi)\) plane.

Thus, to discriminate two jets the angular separation of their axes has to be roughly \(\Delta R \gtrsim 0.2\). If opening angles are smaller, the total energy deposit of the resonance decay products can still be measured, but the substructure, i.e. the energy sharing between the decay products, becomes opaque. To maintain the ability to separate between signal and QCD-induced backgrounds we propose to utilise track-based reconstruction. Trajectories of charged particles as measured in the tracking detectors provide a much better spatial resolution than the reconstructed calorimeter clusters. Recently, a similar approach was advocated for highly boosted electroweak scale resonances [15–18], for which dedicated taggers have been developed.\textsuperscript{2}

In this work, we use track-based reconstruction techniques to evaluate the sensitivity of general purpose detectors at hadron colliders, with characteristics similar to those of ATLAS [19] and CMS [20], in measuring rare Higgs boson decays into light hadronically decaying resonances. Focusing on the High Luminosity LHC (HL-LHC) regime, our analysis assumes a dataset corresponding to an integrated luminosity of 3000 fb\textsuperscript{-1} collected at center-of-mass energy \(\sqrt{s} = 13\) TeV. We consider two production channels for the Higgs boson: inclusive Higgs boson production and Higgs boson production in association with a hard jet of transverse momentum \(p_T > 150\) GeV.

As two benchmark cases for rare Higgs boson decays into light resonances we consider \(h \to Z(\to \ell\ell) + \eta_c\) and \(h \to Z(\to \ell\ell) + A\), where \(A\) is assumed to be an elementary CP-odd scalar of mass 4 GeV which decays mostly hadronically. The presence of two high-pT isolated leptons from the Z boson decay, ensure an efficient trigger strategy for HL-LHC environment. The characteristics of the \(h \to Z(\to \ell\ell) + \eta_c\) benchmark are expected to be representative of similar decays to vector charmonia (e.g. \(h \to Z(\to \ell\ell) + J/\psi\)), due to similarities in their hadronic decay patterns and small mass differences relative to the scale of the jet momenta relevant in the decays of Higgs boson with a mass of 125 GeV.

\textsuperscript{1} Based on a simple sum of the branching fractions for the observed decays of the \(\eta_c\) into stable hadrons.

\textsuperscript{2} http://www IPPP.dur.ac.uk/~mspannow/webipp/HPTTaggers.html.

The event generation is described in Sect. 2, while Sect. 3 is devoted to the details of the reconstruction of the Higgs boson decay products and event selection. The statistical analysis and expected sensitivity are given in Sect. 4. In Sect. 5 the expected results are interpreted in terms of 2HDM models. We offer a summary of our findings in Sect. 6.

\section{Event generation}

For the simulation of the background contributions we employ a modified version of Sherpa 2.2.0 [21] that was adapted in such a way as to facilitate the simulation of Higgs decays into composite resonances. Parton shower effects, hadronisation, as well as underlying event contributions are taken into account throughout. Both Higgs boson production processes, \(h +\text{jet}\) and inclusive \(h\), are calculated at NLO and matched to the parton shower. Finite top quark mass effects in the gluon fusion production mechanism are taken into account as described in Ref. [22]. The Higgs boson decays \(h \to Z \eta_c\), \(h \to Z A\) as well as the subsequent decay of the pseudoscalar and the \(Z\) boson are calculated perturbatively at leading order using the algorithm and methods described in Ref. [23]. Spin-correlations are thus retained in all resonance decays. The UFO model format, supported by Sherpa, was used for the implementation of an elementary pseudoscalar and its interactions [23,24].

The \(Z + \text{jets}\) production is expected to represent the dominant background in this search with other contributions such as \(t\bar{t}\) production being suppressed to a negligible level by requiring an opposite-charge same-flavour dilepton with an invariant mass consistent with the \(Z\) boson mass. For inclusive \(Z\) boson production (\(Z + \text{jets}\)), we take into account the full dilepton final state in the matrix elements and calculate the core process at NLO. We account for additional hard jet emissions by means of multijet merging techniques [25] and include leading order matrix elements with up to two additional jets in the setup.

We process the generated events with the DELPHES\textsuperscript{2} fast simulation framework [26], which uses parametrised descriptions of the response of particle physics detectors to provide reconstructed physics objects, allowing a realistic data analyses to be performed. As an example of a general purpose LHC detector, the default ATLAS configuration card included in DELPHES is used.

\section{Reconstruction setup and selection}

\subsection{Leptonic \(Z\) boson decay reconstruction}

The reconstruction of \(Z \to \ell\ell\) decays begins with the identification of isolated lepton (electron or muon) candidates. Reconstructed leptons are required to satisfy \(p_T > 8\) GeV.
3.2 Hadronic resonance reconstruction

The reconstruction of hadronically decaying resonances within events begins with a search for anti-\(k_t\) calorimeter jets with \(R = 0.4\), seeded by clusters of calorimeter energy deposits. Calorimeter jets are required to have \(p_T > 30\) GeV and \(|\eta| < 2.5\). Any jets which are within \(\Delta R < 0.3\) of leptons forming a \(Z \rightarrow \ell\ell\) candidate are rejected. Following the identification of such a jet, the jet constituents are used to seed a search for an anti-\(k_t\) calorimeter jet with \(R = 0.2\). The identification of an \(R = 0.2\) jet from the constituents of the initial \(R = 0.4\) jet is required to be successful. This procedure, i.e. the reconstruction of anti-\(k_t\) \(R = 0.2\) jets from the constituents of identified \(R = 0.4\) jets, is repeated for track jets, seeded by reconstructed charged particles. Track jets are associated to calorimeter jets by a simple spatial matching, based on a requirement of \(\Delta R < 0.4\) between the axes of the \(R = 0.4\) calorimeter and track jets. Only jets reconstructed with both calorimeter and track components are considered for further analysis and at least one such jet is required to be reconstructed.

To distinguish hadronically decaying charmonium states or light scalars from the copious production of low \(p_T\) jets, a boosted decision tree (BDT) is used through the TMVA package [27]. The following variables are used as input to the BDT:

- The mass of the \(R = 0.4\) track and calorimeter jets, as the jets in the signal are expected to be close to the mass of the light resonance.
- The number of track constituents associated with the \(R = 0.4\) and \(R = 0.2\) track jets, as the signal is expected to have a lower track multiplicity given the upper bound imposed by the light resonance mass.
- The ratio of the \(R = 0.2\) calorimeter (track) jets \(p_T\) to the \(p_T\) of the associated \(R = 0.4\) calorimeter (track) jet, this quantity is expected to prefer values more toward unity in the signal case where a narrow boosted topology is expected, a wider distribution expected from the QCD jet background.
- The spatial separation, \(\Delta R\), between the leading \(p_T\) track within the \(R = 0.4\) track jet and the jet axis.
- The ratio of the highest track \(p_T\) to the \(p_T\) of the \(R = 0.4\) track jet.

The final variables are designed to exploit the fact that in the signal we find on average fewer charged tracks and, due to the very small resonance mass, a smaller active area of the jet.

The performance of the BDT is summarised in Fig. 1, where the background rejection is shown as a function of the signal efficiency. Higgs decays into a composite light resonance \(\eta_c\) and Higgs decays into an elementary pseudoscalar \(A\), which in turn decays hadronically, are considered separately. For the elementary pseudoscalar, individual curves for the case in which it decays into a pair of quarks (\(c\bar{c}\) taken as an example) and for the case in which in decays into a pair of gluons are shown. These pseudoscalar decay modes will be of relevance for the interpretation of our results in the context of 2HDMs in Sect. 5. Examples of the distributions of the variables used to train the BDT are shown in Fig. 2. The most important variables in terms of discrimination between signal and background are found to be the jet masses, followed by the number of track constituents associated with the track jets.
3.3 Selection of $h \rightarrow ZA$ and $h \rightarrow Z \eta_c$ decays

Events containing at least one hadronic decay candidate and one $Z \rightarrow \ell\ell$ candidate are considered for further analysis. In the case of the $h + \text{jet}$ production channel, an additional $R = 0.4$ anti-$k_t$ calorimeter jet with $p_T > 150$ GeV and $|\eta| < 2.5$ is required (no substructure or matching track jet is required). The single $Z$ boson candidate with $m_{\ell\ell}$ closest to the $Z$ boson mass is chosen to form the $h \rightarrow ZA(\eta_c)$ candidate. If multiple hadronic decay candidates are reconstructed, the candidate which when paired with the $Z \rightarrow \ell\ell$ candidate has an invariant mass closest to $m_h = 125$ GeV is chosen. Finally, the transverse momentum of the $h$ candidate is required to exceed 20 GeV. The invariant mass of the jet–dilepton system is shown for the inclusive and $h + \text{jet}$ production channels in Fig. 3.

The BDT response is shown for both the signal and the background contributions to the inclusive and $h + \text{jet}$ production channels in Fig. 4.

4 Statistical analysis and results

The expected performance of the analysis is used to evaluate expected 95 % CL limits on the branching fractions $BR(h \rightarrow ZA)$, in the cases where $BR(A \rightarrow gg) = 1.0$ or $BR(A \rightarrow c\bar{c}) = 1.0$, and $BR(h \rightarrow Z\eta_c)$. The yields of signal and background events within $110 \text{ GeV} < m_{\ell\ell} < 140 \text{ GeV}$ are used to evaluate the limits. To exploit the additional sensitivity offered by the BDT, a requirement on the BDT response is imposed. The value of this requirement is optimised to provide the best limit on the branching fractions of interest. The expected 95 % CL limits on the branching fractions of interest are shown Table 1. Branching fraction limits at the 1 % level can be expected. The inclusive production channel is found to be slightly more sensitive than the $h + \text{jet}$ channel.

In addition to the channels described, Higgs boson production in association with a leptonically decaying $Z$ boson was also considered as a possible channel to gain additional sensitivity. Initial studies into this channel demonstrated improved signal-to-background ratios when compared to the two channels constituting the main study, though the substantially lower number of signal events resulted in expected branching fraction limits that were up to an order of magnitude higher than the inclusive and $h + \text{jet}$ channels.

5 Constraints on the 2HDM parameter space

With a focus on the HL-LHC, we assume the Higgs boson couplings to be tightly constrained to SM-like values. Assuming no evidence for new physics in the HL-LHC data, any 2HDM scenario compatible with the observations would therefore necessarily be close to the alignment limit. It has been pointed out in Ref. [5] that a light pseudoscalar $A$ with mass below 10 GeV can be accommodated in this limit, particularly in Type I models, which we consider here. A pseudoscalar that light can decay into pairs of fermions through tree-level interactions or into pairs of gluons and photons through loop-induced couplings. In Type I models, the tree-level couplings to fermions are essentially given by the
Fig. 3 The invariant mass distribution of the jet–dilepton system (with no BDT based selection applied) in inclusive $h$ production (left) and $h +$ jet production (right) is shown for $A \rightarrow gg$ (top), $A \rightarrow c\bar{c}$ (middle) and $\eta_c \rightarrow$ hadrons (bottom) signals in comparison to the background contribution. The signal contribution is multiplied by ten to improve visibility.

Fermion masses times a universal factor of $\cot(\beta)$. A considerable hadronic branching fraction hence arises from decays into quark pairs, gluon pairs, or indirectly from decays into pairs of tau leptons that decay into hadrons subsequently. As shown in Fig. 1, the performance of our analysis is fairly insensitive to the details of the hadronic decay mode of the pseudoscalar. The results of our analysis can therefore directly be used in order to constrain such models. To the best of our knowledge, no detailed analysis of this final state has been provided in the literature so far.

In order to assess the constraining power of our results, we perform a parameter scan for a fixed benchmark pseudoscalar mass of $m_A=4\,\text{GeV}$. For the chosen benchmark value of $m_A$, decays into tau leptons and charm quarks dominate. Decays into gluon pairs contribute a branching fraction at the percent level. Overall, we obtain $\text{BR}(A \rightarrow \text{hadrons}) \approx 82\%$. 

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In our parameter scan, we calculate the branching ratio relevant for the interpretation of our results, $BR(h \to ZA)$, for each parameter point. The corresponding partial decay width is given by

$$\Gamma(h \to ZA) = \frac{|p|}{8\pi m_h^2} |\mathcal{M}(h \to ZA)|^2 = \frac{g_{hZA}^2}{2\pi} \frac{|\vec{p}|^3}{m_Z^2},$$

at tree level, where $\vec{p}$ is the three-momentum of either of the two decay products in the rest frame of the Higgs boson. The $hZA$-coupling is given by

$$g_{hZA} = \frac{e \cos(\beta - \alpha)}{2 \cos \theta_W \sin \theta_W}. \quad (5.2)$$

The partial decay width $\Gamma(h \to ZA)$ therefore vanishes in the strict alignment limit with $\cos(\beta - \alpha) = 0$. The corresponding branching fraction, however, becomes sizable already for small $\cos(\beta - \alpha)$ if the decay $h \to AA$ does not contribute substantially to the Higgs boson total width. We therefore focus on the parameter region, where $g_{hZA}^{SM} = 0$ at tree level, which implies [28]

$$m_{12}^2 = (2m_h^2 + m_Z^2) \sin(2\beta)/4.0. \quad (5.3)$$
Table 1 The expected 95 % CL limits on the branching fractions of interest for both the inclusive and the $h + \text{j}et$ channels, assuming $3000\,\text{fb}^{-1}$ at $\sqrt{s} = 13\,\text{TeV}$

| Channel | $BR$ 95 % confidence level upper limit |
|---------|--------------------------------------|
| $h \to ZA(\to gg)$ | 2.0 % | 2.1 % | 2.0 % |
| $h \to ZA(\to c\bar{c})$ | 3.5 % | 3.9 % | 3.7 % |

To ensure alignment, we perform a uniform scan with $\sin(\beta - \alpha) \in [0.99, 1.0]$. In this regime, we can assume the production cross sections of the 125 GeV Higgs to be SM-like and directly apply our previously obtained limit on $BR(h \to ZA)$. Note, however, that the limit must be applied to $BR(h \to ZA) \times BR(A \to \text{hadrons})$, since $BR(A \to \text{hadrons}) = 1$ was assumed previously. The remaining free parameters of the model are uniformly varied in the intervals $m_H \in [130, 600]$ GeV, $m_{H^\pm} \in [50, 600]$ GeV, and $\tan \beta \in [0.1, 5.0]$. We calculate the physical spectrum and the relevant branching fractions with 2HDMC version 1.7.0 [29].

For each point we check for vacuum stability of the potential, tree-level unitarity using the corresponding functional dependence of these quantities given by Eq. (5.1), assuming for simplicity $\Gamma_{\text{tot}} = \Gamma_{h/\text{tot}}^h$. For large $\cos(\beta - \alpha)$, this assumption is violated due to the opening of further decay channels. At small $\cos(\beta - \alpha)$, however, the corresponding approximation proves to be reasonable for parameter points that pass the applied phenomenological constraints. As illustrated in Fig. 5, the scanned parameter space can effectively be constrained to very small values of $\cos^2(\beta - \alpha)$ by applying our expected limit on $BR(h \to ZA)$. In fact, we find that no parameter point with $\cos^2(\beta - \alpha) > 0.0035$ survives the limit set by the analysis presented above, translating to $\sin(\beta - \alpha) \gtrsim 0.998$ in the scanned subspace of parameters. Correspondingly, a mere 12 % of the parameter points displayed in Fig. 5 fall in the region of allowed values for $BR(h \to ZA) \times BR(A \to \text{hadrons})$ after applying the limit on $BR(h \to ZA)$ obtained above.

6 Summary

Searches for rare and exclusive Higgs boson decays are at the core of the program of the High Luminosity LHC. The observation of Higgs boson decays into light elementary or composite resonances would be evidence for the existence of physics beyond the Standard Model.

While previous experimental strategies to reconstruct light resonances relied entirely on their leptonic decay products, in this work, we evaluated the prospects for their discovery in the often dominant hadronic decay channels. We have focused on the Higgs boson production processes with the largest cross sections, $pp \to h$ and $pp \to h + \text{j}et$, with subsequent decays $h \to ZA$ or $h \to Z \eta_c$. The former is present in many multi-Higgs extensions of the Standard Model, while observing the latter at a branching ratio of $BR(h \to Z \eta_c) \geq 10^{-3}$ could indicate an enhanced Higgs-charm coupling.

The decay products of light resonances with masses below a few GeV that arise from Higgs decays are highly collimated, i.e. they get emitted into a small area of the detector. In such scenarios jet substructure is an indispensable tool to retain sensitivity in discriminating signal from large QCD-induced backgrounds. In particular, by exploiting the

![Fig. 5 Distribution of scanned parameter points in the $\cos^2(\beta - \alpha)$ vs. $BR(h \to ZA) \times BR(A \to \text{hadrons})$ plane. The color-coding denotes the density of points in the respective areas as indicated by the color bar. We also display the tree-level functional relationship between $\cos^2(\beta - \alpha)$ and $BR(h \to ZA) \times BR(A \to \text{hadrons})$, assuming $\Gamma_{h/\text{tot}}^h = \Gamma_{\text{tot}}^{h/\text{tot}}$. The dashed line shows the expected 95 % CL upper limit on the displayed branching fraction. All points above this line are expected to be excluded by the analysis presented here.](image)
improved angular resolution of track-based observables, a good signal-to-background discrimination can be achieved, which results in a limit on the branching ratios of $O(1)$ % for a data sample corresponding to $3000 \text{fb}^{-1}$ at $\sqrt{s} = 13 \text{TeV}$.

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References

1. ATLAS collaboration, Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC. Phys. Lett. B 716, 1–29 (2012). arXiv:1207.7214

2. CMS collaboration, Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC. Phys. Lett. B 716, 30–61 (2012). arXiv:1207.7235

3. ATLAS, CMS collaboration, Measurements of the Higgs boson production and decay rates and constraints on its couplings from a combined ATLAS and CMS analysis of the LHC pp collision data at $\sqrt{s} = 7$ and 8 TeV. arXiv:1606.02266

4. D. Curtin et al., Exotic decays of the 125 GeV Higgs boson. Phys. Rev. D 90, 075004 (2014). arXiv:1312.4992

5. J. Bernet, J.F. Gunion, H.E. Haber, Y. Jiang, S. Kraml, Scrutinizing the alignment limit in two-Higgs-doublet models: $m_h = 125$ GeV. Phys. Rev. D 92, 075004 (2015). arXiv:1507.00933

6. D. Curtin, R. Essig, S. Gori, J. Shelton, Illuminating dark photons with high-energy colliders. JHEP 02, 157 (2015). arXiv:1412.0018

7. G.T. Bodwin, F. Petriello, S. Stoynev, M. Velasco, Higgs boson decays to quarkonia and the $H\ell\gamma$ coupling. Phys. Rev. D 88, 053003 (2013). arXiv:1306.5770

8. A.L. Kagan, G. Perez, F. Petriello, Y. Soreq, S. Stoynev, J. Zupan, Exclusive window onto Higgs Yukawa couplings. Phys. Rev. Lett. 114, 101802 (2015). arXiv:1406.1722

9. G. Isidori, A.V. Manohar, M. Trott, Probing the nature of the Higgs-like Boson via $h \rightarrow V\ell\ell$ decays. Phys. Lett. B 728, 131–135 (2014). arXiv:1305.0663

10. M. König, M. Neubert, Exclusive radiative Higgs decays as probes of light-Quark Yukawa couplings. JHEP 08, 012 (2015). arXiv:1505.03870

11. ATLAS collaboration, Search for Higgs and Z Boson decays to $J/\psi$ and $\Upsilon(nS)$ with the ATLAS Ddetector. Phys. Rev. Lett. 114, 121801 (2015). arXiv:1501.03276

12. CMS collaboration, Search for a Higgs boson decaying into $\gamma\gamma \rightarrow \ell\ell\gamma$ with low dilepton mass in pp collisions at $\sqrt{s} = 8 \text{TeV}$. Phys. Lett. B 753, 341-362 (2016). arXiv:1507.03031

13. ATLAS collaboration, Search for Higgs and Z boson decays to $\phi\gamma$ with the ATLAS detector. Phys. Rev. Lett. 117(11), 111802 (2016). arXiv:1607.03400

14. Particle Data Group collaboration, K. A. Olive et al., Review of particle physics. Chin. Phys. C 38, 090001 (2014)

15. A. Katz, M. Son, B. Tweedie, Jet substructure and the search for neutral spin-one resonances in electroweak boson channels. JHEP 03, 011 (2011). arXiv:1010.5253

16. S. Schätzle, M. Spannowsky, Tagging highly boosted top quarks. Phys. Rev. D 89, 014007 (2014). arXiv:1308.0540

17. A.J. Larkoski, F. Maltoni, M. Selvaggi, Tracking down hyper-boosted top quarks. JHEP 06, 032 (2015). arXiv:1503.03347

18. M. Spannowsky, M. Stoll, Tracking new physics at the LHC and beyond. Phys. Rev. D 92, 054033 (2015). arXiv:1505.01921

19. ATLAS collaboration, The ATLAS experiment at the CERN large hadron collider. JINST 3, S08003 (2008)

20. CMS collaboration, The CMS experiment at the CERN LHC. JINST 3, S08004 (2008)

21. T. Gleisberg, S. Höche, F. Krauss, M. Schönherr, S. Schumann, F. Siegert et al., Event generation with SHERPA 1.1. JHEP 02, 007 (2009). arXiv:0811.4622

22. M. Buschmann, D. Goncalves, S. Kuttimalai, M. Schönherr, F. Krauss, T. Plehn, Mass effects in the Higgs-Gluon coupling: boosted vs off-shell production. JHEP 02, 038 (2015). arXiv:1410.5806

23. S. Höche, S. Kuttimalai, S. Schumann, F. Siegert, Beyond standard model calculations with Sherpa. Eur. Phys. J. C 75, 135 (2015). arXiv:1412.6478

24. C. Degrande, C. Duhr, B. Fuks, D. Grellscheid, O. Mattelaer, T. Reiter, UFO—the Universal FeynRules output. Comput. Phys. Commun. 183, 1201–1214 (2012). arXiv:1108.2040

25. S. Höche, F. Krauss, M. Schönherr, F. Siegert, QCD matrix elements + parton showers: the NLO case. JHEP 04, 027 (2013). arXiv:1207.5030

26. DELPHES 3 collaboration, J. de Favereau, C. Delaere, P. Demin, A. Giammanco, V. Lemaitre, A. Mertens et al., DELPHES 3, a modular framework for fast simulation of a generic collider experiment. JHEP 02, 057 (2014). arXiv:1307.6346

27. A. Hoecker et al., TMVA—toolkit for multivariate data analysis. PoS ACAT, 040 (2007). arXiv:physics/0703039

28. M. Casolino, T. Farooque, A. Juste, T. Liu, M. Spannowsky, Probing a light CP-odd scalar in di-top-associated production at the LHC. Eur. Phys. J. C 75, 498 (2015). arXiv:1507.07004

29. D. Eriksson, J. Rathsman, O. Stål, 2HDMC: two-Higgs-Doublet model calculator physics and manual. Comput. Phys. Commun. 181, 189–205 (2010). arXiv:0902.0851

30. M.E. Peskin, T. Takeuchi, A new constraint on a strongly interacting Higgs sector. Phys. Rev. Lett. 65, 964–967 (1990)

31. M.E. Peskin, T. Takeuchi, Estimation of oblique electroweak corrections. Phys. Rev. D 46, 381–409 (1992)

32. Glüker Group collaboration, M. Baak et al., The global electroweak fit at NNLO and prospects for the LHC and ILC. Eur. Phys. J. C 74, 3046 (2014). arXiv:1407.3792

33. P. Bechtle, O. Brein, S. Heinemeyer, G. Weiglein, K.E. Williams, HiggsBounds: confronting arbitrary Higgs sectors with exclusion bounds from LEP and the Tevatron. Comput. Phys. Commun. 181, 138–167 (2010). arXiv:0811.4169

34. P. Bechtle, O. Brein, S. Heinemeyer, G. Weiglein, K.E. Williams, HiggsBounds 2.0.0: confronting neutral and charged Higgs sector predictions with exclusion bounds from LEP and the Tevatron. Comput. Phys. Commun. 182, 2605–2631 (2011). arXiv:1102.1898

35. P. Bechtle et al., Recent developments in HiggsBounds and a preview of HiggsSignals. PoS CHARGED2012, 024 (2012). arXiv:1301.2345

36. P. Bechtle, S. Heinemeyer, O. Stål, T. Stefaniak, G. Weiglein, Applying exclusion likelihoods from LHC searches to extended Higgs sectors. Eur. Phys. J. C 75, 421 (2015). arXiv:1507.06706
37. P. Bechtle, S. Heinemeyer, O. Stål, T. Stefaniak, G. Weiglein, *HiggsSignals*: confronting arbitrary Higgs sectors with measurements at the Tevatron and the LHC. Eur. Phys. J. C **74**, 2711 (2014). arXiv:1305.1933

38. P. Bechtle, S. Heinemeyer, O. Stål, T. Stefaniak, G. Weiglein, Probing the standard model with Higgs signal rates from the Tevatron, the LHC and a future ILC. JHEP **11**, 039 (2014). arXiv:1403.1582