Interplay of superconductivity and magnetism in FeSe\textsubscript{1-x}Te\textsubscript{x} compounds.
Pressure effects.

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The influence of uniform pressures \( P \) up to 5 kbar on the superconducting transition temperature \( T_c \) was studied for the FeSe\textsubscript{1-x}Te\textsubscript{x} \((x = 0, 0.85, 0.88 \) and 0.9\) system. For the first time, we observed a change in sign of the pressure effect on \( T_c \) when going from FeSe to tellurium rich alloys. This has allowed to specify the pressure derivative \( dT_c/dP \) for the system as a function of composition. The observed dependence was compared with results of the \textit{ab initio} calculations of electronic structure and magnetism of FeSe, FeTe and FeSe\textsubscript{0.5}Te\textsubscript{0.5}, and also with our recent experimental data on pressure effects on magnetic susceptibilities of FeSe and FeTe compounds in the normal state. This comparison demonstrates a competing interplay between superconductivity and magnetism in tellurium rich FeSe\textsubscript{1-x}Te\textsubscript{x} compounds.

Keywords: Fe-based superconductors, FeSe\textsubscript{1-x}Te\textsubscript{x}, electronic structure, magnetic susceptibility, pressure effects

I. INTRODUCTION

For the most families of recently discovered class of the Fe-based high-temperature superconductors (HTSC) the emergence of superconductivity with doping or under uniform pressure is accompanied by suppression of the magnetic ordering \[1-4\]. It is widely believed then that spin fluctuations play an important role in formation of the Cooper pairs \[5-7\]. Nevertheless, as shown e.g. in Ref. \[8\], for many Fe-based HTSCs the experimental values of superconducting transition temperatures are well described in the framework of the electron-phonon mechanism of pairing. The close interrelation of magnetism and superconductivity determines the importance of further studying of magnetic and superconducting properties and their evolution under variations of composition, pressure, etc. for understanding HTSC mechanism in the considered new class of iron compounds. One of representatives of this class is the system of FeSe\textsubscript{1-x}Te\textsubscript{x} chalcogenides, which possesses the simplest crystal structure among iron-based superconductors, that favors to experimental and theoretical studying the effects of chemical substitution and high pressures on its properties.

Superconducting properties of FeSe\textsubscript{1-x}Te\textsubscript{x} are characterized by nonmonotonic dependence of transition temperatures \( T_c \) on composition. There is a noticeable growth from \( T_c \approx 8 \) K for \( x = 0 \) to the maximum value \( \sim 15 \) K at \( x \approx 0.5 \) with the subsequent falling to 0 K near \( x \approx 0.9 \) (see, for example, Ref. \[9\] and references therein). Also, in FeSe compound the extremely large rise of \( T_c \) up to 35–37 K takes place with pressure \( P \approx 70–80 \) kbar \[10,11\]. The similar behavior of \( T_c \) under pressure was also observed in FeSe\textsubscript{0.5}Te\textsubscript{0.5} compound \[12,13\]. With further increase of \( x \) in FeSe\textsubscript{1-x}Te\textsubscript{x} a tendency to reduction of the positive pressure effect is expected with even probable change of its sign, as it was observed in the related tellurium rich FeSe\textsubscript{0.2}Te\textsubscript{0.8} alloy \[9\]. This alleged change in sign of the pressure effect on \( T_c \) in FeSe\textsubscript{1-x}Te\textsubscript{x} under substitution of Te for Se could also explain the reason of unsuccessful attempts to observe superconductivity in FeTe under pressures up to 190 kbar \[14,15\].

Magnetic properties of FeSe\textsubscript{1-x}Te\textsubscript{x} system were investigated in a number of works \[16-22\], however, data on the magnetic susceptibility in the normal state remain incomplete and quantitatively inconsistent. This is caused not only by a different quality of the samples used, but also by the existence in them of impurities of iron and its secondary magnetic phases which considerably mask their intrinsic magnetic susceptibility and must be carefully taken into account \[22\]. The most adequate experimental data indicate that the susceptibility of FeSe\textsubscript{1-x}Te\textsubscript{x} compounds increases gradually with
Te content, being in FeTe about one order of magnitude lager than that of FeSe. Moreover, FeTe compound becomes magnetically unstable, and the antiferromagnetic ordering has been observed at temperatures about 70 K (see e.g. Ref. [17]).

It should be noted that the largest rise of magnetic susceptibility in the normal state, $\chi(x)$, with increase of $x$ is observed in tellurium rich compounds, where, in turn, the $T_c(x)$ dependence falls steeply down and FeTe compound is not superconductor under ambient conditions. This allows to assume a competing interplay between magnetism and superconductivity, at least for this range of compositions. In order to shed more light on the relationship between magnetic and superconducting properties in FeSe$_{1-x}$Te$_x$ system, it is very important to study evolution of these properties under high pressure. For this purpose in the present work we investigated the influence of hydrostatic pressure on the superconducting transition temperature, mainly in tellurium rich FeTe(Se) compounds. The obtained experimental results were compared with available data on behavior of magnetic susceptibility under pressure for the basic compounds FeSe [25] and FeTe [14, 26], also supplemented by calculated pressure dependencies of electronic structure and magnetic susceptibility for FeSe$_{0.5}$Te$_{0.5}$ compound.

II. EXPERIMENTAL DETAILS AND RESULTS

The single crystals of FeSe$_{0.96}$ superconductor (hereinafter referred to as FeSe) were grown during 50 days in evacuated quartz ampoules using the AlCl$_3$/KCl flux technique with a constant temperature gradient along the ampoule length [27]. Temperature of the hot end of the ampoule was kept at 427°C, when its more cold end was at about 380°C. A similar method was employed for the synthesis of tellurium-rich single crystals of Fe$_{1+\delta}$Se$_{1-x}$Te$_x$ superconductors ($\delta \sim 0.05$, $x = 0.85$, 0.88 and 0.90). In this case we used the KCl/NaCl salt mixture and temperatures of the hot and cold ends of the ampoule were 750°C and about 700°C, respectively. The duration of the synthesis was 20-25 days. Typical dimensions of the produced plate-like single crystals were (1 - 3) $\times$ (1 - 3) $\times$ (0.2 - 0.3) mm$^3$. Their tetragonal $P4/nnm$ structure was demonstrated at room temperature by an x-ray diffraction technique. The crystals composition was determined using energy dispersive X-ray spectroscopy, performed on a CAMECA SX100 (15 keV) analytical scanning electron microscope, with an accuracy of the components ratio not worse than 2% (for details, see [24, 27]).

The measurements of magnetic properties were performed using a SQUID magnetometer (MPMS-XL5 Quantum Design) equipped with a miniature high-pressure cell of a piston-cylinder type (similar to that described in Ref. [28]). The cell was made of non-magnetic CuBe alloy with the inside and outside diameters of 1.6 mm and 5 mm, respectively. Polyethylsiloxane liquid PES-3 was used as a hydrostatic pressure-transmitting medium. The value of pressure at low temperatures was determined according to the known pressure dependence of the superconducting transition temperature for a sample of pure tin [29], located inside the cell close to the measured sample. The corresponding error did not exceed 0.2 kbar.

Fig. 1 shows the temperature dependencies of magnetic moment $M(T)$ for FeSe at different values of pressure, which were measured under cooling of the sample in zero magnetic field (ZFC) followed by its heating in the field $H = 10$ Oe. Resulted from Fig. 1 pressure dependence of the superconducting transition temperature $T_c$, determined from here on by the onset of the transition, is given in Fig. 2. Within the experimental errors and the operating range of pressure, this dependence appeared to be close to linear that allows to evaluate the pressure derivative $dT_c/dP$.

The $M(T)$ dependencies for tellurium-rich FeSe$_{1-x}$Te$_x$ compounds were measured at different pressures in ZFC regime, and are shown in Fig. 3. They demonstrate clearly defined negative pressure effect on the superconducting transition temperature. Experimental values of $T_c$ and its pressure derivative for all investigated samples are listed in Table 1. As is evident from the presented data, the pressure effects on $T_c$ in the tellurium rich FeSe$_{1-x}$Te$_x$ compounds are comparable in magnitude with that for FeSe but have opposite negative sign.

![Figure 1: Temperature dependence of the magnetic moment of FeSe, measured in magnetic field $H = 10$ Oe at different pressures.](image-url)
III. CALCULATIONS OF ELECTRONIC STRUCTURE AND MAGNETIC SUSCEPTIBILITY OF FeSe$_{0.5}$Te$_{0.5}$ COMPOUND

For calculations of electronic structure of FeSe$_{0.5}$Te$_{0.5}$ compound we employed the relativistic full potential LMTO method (FP-LMTO, RSPt implementation [30, 31]). The exchange-correlation potential was treated within the local density approximation (LDA [32]) of the density functional theory (DFT). The calculations were carried out for a supercell $2 \times 2 \times 1$, constructed by double translations of the unit cell for the ordered tetragonal phase of FeSe and FeTe along the crystallographic [100] and [010] directions, by using experimental values of crystal lattice parameters for FeSe$_{0.5}$Te$_{0.5}$ from Refs. [33–35]. The calculated density of electronic states (DOS) $N(E)$ of the paramagnetic FeSe$_{0.5}$Te$_{0.5}$ compound is presented in Fig. 4. The Fermi level $E_F$ is situated in the region of a local flat plateau of $N(E)$, where the main contribution to DOS comes from the $d$-states of iron. Such position of $E_F$ implies a weak temperature dependence of the spin susceptibility in FeSe$_{0.5}$Te$_{0.5}$, which is consistent with available experimental data for this compound [16, 22, 23].

To evaluate the paramagnetic susceptibility of

Table I: Superconducting transition temperature $T_c$ and its pressure derivative $dT_c/dP$ for FeSe$_{1-x}$Te$_x$ compounds.

| Composition | $T_c$ (K) | $dT_c/dP$ (K/kbar) |
|-------------|----------|-------------------|
| $x=0$       | 9.12     | 0.78 ± 0.05       |
| $x=0.85$    | 11.62    | −0.31 ± 0.05      |
| $x=0.88$    | 11.05    | −0.40 ± 0.05      |
| $x=0.90$    | 9.71     | −0.40 ± 0.1       |

Figure 2: Pressure dependence of the superconducting transition temperature for FeSe.

Figure 3: Temperature dependencies of the magnetic moment measured in $H=10$ Oe at two values of pressure for tellurium-rich FeSe$_{1-x}$Te$_x$ compounds: (a) - $x=0.85$, (b) - $x=0.88$, (c) - $x=0.9$. Arrows 1 and 2 denote $T_c$ at zero and finite values of pressure, respectively.

Figure 4: Density of electronic states $N(E)$ of FeSe$_{0.5}$Te$_{0.5}$ compound. The Fermi level position at 0 eV is marked by a vertical line.
The FP-LMTO calculations of field-induced spin and orbital (Van Vleck) magnetic moments were carried out with the approach described in Ref. 31 within the local spin density approximation (LSDA) of DFT. The relativistic effects, including spin-orbit coupling, were incorporated, and the effect of an external magnetic field \( \mathbf{B} \) was taken into account self-consistently by means of the Zeeman term:

\[
\mathcal{H} = \mu_B \mathbf{B} \cdot (2\mathbf{s} + \mathbf{l}),
\]

(1)

Here \( \mu_B \) is the Bohr magneton, \( \mathbf{s} \) and \( \mathbf{l} \) are the spin and orbital angular momentum operators, respectively. The ratio of the field-induced magnetizations to the field strength (\( \mathbf{B} = 10 \) T) provided corresponding spin and orbital components of magnetic susceptibilities, \( \chi_{\text{spin}} \) and \( \chi_{\text{orb}} \), respectively.

According to results of the calculations, the exchange-enhanced spin paramagnetism \( \chi_{\text{spin}} \) appears to be the main contribution to magnetic susceptibility of FeSe\(_{0.5}\)Te\(_{0.5}\) compound. Within the Stoner model, this contribution can be presented as:

\[
\chi_{\text{spin}} = S \mu_B^2 N(E_F),
\]

where \( S \) is the Stoner factor, \( N(E_F) \) DOS at the Fermi level, \( \mu_B \) the Bohr magneton. Using the calculated values of spin magnetic susceptibility of FeSe\(_{0.5}\)Te\(_{0.5}\) compound, \( \chi_{\text{spin}} \approx 0.6 \times 10^{-5} \text{ emu/mol} \), and DOS at the Fermi level, \( N(E_F) \approx 1.85 \text{ eV}^{-1} \), we have obtained the estimation of the Stoner factor: \( S \approx 10 \). It should be noted that the above listed calculated value of \( \chi_{\text{spin}} \) is in agreement with the experimental magnetic susceptibility of FeSe\(_{0.5}\)Te\(_{0.5}\) compound in the normal state (see Refs. 16, 23). This confirms the dominating role of the spin contribution to magnetism of FeSe\(_{0.5}\)Te\(_{0.5}\) compound, that is, apparently, characteristic for the whole FeSe\(_{1-x}\)Te\(_x\) system 24, 25, 26.

By using the experimental data of Ref. 35 on evaluation of the lattice parameters of FeSe\(_{0.5}\)Te\(_{0.5}\) under uniform compression, we calculated the behavior of density of electronic states at the Fermi level. For the region of small pressures (0–10 kbar) we established the growth of \( N(E_F) \) with the rate of \( \partial \ln N(E_F)/\partial P \approx 1 \text{ Mbar}^{-1} \). We should note that such behavior of \( N(E_F) \) correlates with increase of the superconducting transition temperature in FeSe\(_{0.5}\)Te\(_{0.5}\) under pressure 12, 13.

Within the considered above method of calculation of magnetic susceptibility, we also investigated the dependence of \( \chi \) in FeSe\(_{0.5}\)Te\(_{0.5}\) compound on the uniform pressure. By direct calculations of the field-induced magnetic moments, we have obtained the value of pressure derivative of paramagnetic susceptibility, \( d \ln \chi / dP \approx 13 \text{ Mbar}^{-1} \), which appeared to be close to the corresponding values in FeSe and FeTe (see Tab. 1). In order to clarify the mechanism of the strong increase of magnetic susceptibility in FeSe\(_{0.5}\)Te\(_{0.5}\) under pressure, we calculated value of \( \chi \) as a function of the unit cell volume \( V \) and the internal structural parameter \( Z \), which determines the relative height of chalcogen atoms over the plane of iron atoms. Then the corresponding pressure effect on \( \chi \) can be presented as follows:

\[
\frac{d \ln \chi}{dP} = \frac{\partial \ln \chi}{\partial \ln V} \frac{d \ln V}{dP} + \frac{\partial \ln \chi}{\partial Z} \frac{dZ}{dP},
\]

(2)

By small variations of the cell volume \( V \) and the structural parameter \( Z \) near their experimental values, the following partial derivatives of paramagnetic susceptibility for FeSe\(_{0.5}\)Te\(_{0.5}\) were calculated to be \( \partial \ln \chi / \partial \ln V \approx 10 \) and \( \partial \ln \chi / \partial Z \approx 90 \). The necessary values for the compressibility of FeSe\(_{0.5}\)Te\(_{0.5}\), \( d \ln V / dP = -3.1 \text{ Mbar}^{-1} \), and behavior of parameter \( Z \) under pressure, \( dZ / dP \approx 0.49 \text{ Mbar}^{-1} \), were taken from Ref. 36. By substituting these values of these parameters in Eq. (2) we have found that the calculated in this work large positive pressure effect on \( \chi \) in FeSe\(_{0.5}\)Te\(_{0.5}\) is related to the strong sensitivity of susceptibility to the parameter \( Z \) and its change under pressure, that determines the dominant positive contribution.

\[\text{IV. DISCUSSION}\]

Experimental values of superconducting transition temperatures for the investigated in this work compounds are in agreement with the literature data (see Fig. 5). The most studied range of compositions (\( x \geq 0.4 \)) is characterized by the sharp reduction of \( T_c \) with increasing \( x \).
The trend is also consistent with the value of $d\chi/dP$, for $\text{FeSe}_{1-x}\text{Te}_x$ compounds. Experimental temperatures are specified in brackets, results of calculations correspond to $T = 0$ K. The data for FeTe are referred to the paramagnetic state.

| Compound     | $\chi$ | $\chi$ |
|--------------|--------|--------|
| FeSe         | 10 $\pm$ 3 (78 K)$^{a}$ | $\sim 8$ $^{a}$ |
| FeSe$_{0.5}$Te$_{0.5}$ | $\sim 9$ (20 K)$^{b}$ | $\sim 13$ |
| FeTe         | 23 $\pm$ 1.5 (78 K)$^{c}$ | $\sim 20$ $^{c}$ |

$^a$ – from Ref. 29, $^b$ – from Ref. 26, $^c$ – from magnetization data of Ref. 14.

Let us consider now the evolution of superconducting and magnetic properties of $\text{FeSe}_{1-x}\text{Te}_x$ compounds under uniform pressure. Experimental values for pressure derivatives of the superconducting transition temperature are given in Fig. 6, which include the known published data and the results of this work. Apparently, the available data describe the monotonous reduction of the pressure effect in $T_c$ in process of selenium substitution with tellurium, and the change of its sign at $x \sim 0.8$. This trend is also consistent with the value of $dT_c/dP \simeq -0.25$ K/kbar for the related $\text{FeS}_{0.2}\text{Te}_{0.8}$ compound 9.

Unlike the pressure effect on $T_c$, which changes its sign as a function of composition (Fig. 5a), the magnetic susceptibility of FeSe(Se) system in the normal state is characterized by substantial growth under pressure for the whole system. This conclusion follows from available experimental data and theoretical estimates for the basic FeSe 25 and FeTe 11,26 compounds, together with the results of present calculations for pressure dependence of magnetic susceptibility in FeSe$_{0.5}$Te$_{0.5}$ compound.

As can be seen from the values of pressure derivative of susceptibility, $d\chi/dP$, given in Table II for considered FeSe(Se) system the pressure effect not only much exceeds its typical value in the exchange-enhanced itinerant paramagnets 31, but also has the opposite positive sign. This implies an unusual for metallic system possibility of transition to the ferromagnetic state under the influence of experimentally achievable pressures. This is particularly the case of FeTe compound where the pressure effect is the largest. In Ref. 26 from the analysis of temperature dependence of susceptibility for FeTe in the paramagnetic region within the Curie-Weiss law, the values of the paramagnetic Curie temperature and its pressure derivative were evaluated to be $\Theta \sim -340$ K and $d\Theta/dP \sim 7$ K/kbar. Corresponding to them rough estimate of the critical pressure for ferromagnetic transition amounts to 35 kbar. This is in reasonable agreement with results of Ref. 40, where the ferromagnetic state was observed in FeTe for the first time under pressures of $P \geq 20$ kbar.

For convenient comparison of the observed pressure effects in superconducting transition temperatures (Fig. 5b) with pressure effects in magnetic susceptibility, the values of pressure derivatives of susceptibility, $d\chi/dP = \chi \times d\ln\chi/dP$, are presented in Fig. 6 for FeSe, FeSe$_{0.5}$Te$_{0.5}$ and FeTe. To evaluate these derivatives we used the corresponding values of $\chi(T \to 0)$ from Ref. 26 (Fig. 5) and the average values of $d\ln\chi/dP$ from Table II. As can be seen in Fig. 6 the presented composition dependencies of pressure effects in magnetic and superconducting properties of $\text{FeSe}_{1-x}\text{Te}_x$ system are strictly opposite to one another. This fact, along with similar trends in behavior of magnetic susceptibility and $T_c$ as function of composition at ambient pressure (Fig. 5), specifies on antagonistic interrelation of magnetism and superconductivity in $\text{FeSe}_{1-x}\text{Te}_x$ system, which is most pronounced in the tellurium rich compositions at $x \geq 0.7$, and the total disappearance of superconductivity for $x \to 1$. In the same range of compositions the strong growth of magnetic susceptibility in the normal state was observed (Fig. 5b). The obtained strictly opposite tendencies in composition dependencies of superconductivity and magnetism in $\text{FeSe}_{1-x}\text{Te}_x$ system allow to assume that interrelation of these phenomena has competing character, at least for the tellurium rich compounds.

![Figure 6](image_url)

Figure 6: (a) The values of $dT_c/dP$ derivative depending on Te composition $x$ in $\text{FeSe}_{1-x}\text{Te}_x$ compounds ($\Theta$ – this work; $\circ$ - [32]; $\blacksquare$ - [13]; $\blacksquare$ - [34]). (b) dependence of the pressure derivative of magnetic susceptibility in the normal state on Te composition $x$ (see more details in the text).
Conclusions

In this work the negative pressure effect on the superconducting transition temperature of tellurium rich FeSe$_{1-x}$Te$_x$ compounds was observed for the first time. The obtained data allowed to establish an overall picture of the composition dependence for the pressure effect on $T_c$, which monotonously decreases with growth of $x$ and changes its sign at $x \sim 0.8$.

Another feature of FeSe$_{1-x}$Te$_x$ compounds is anomalously large and positive pressure effect on magnetic susceptibility in the normal state for all compositions, which grows with substitution of tellurium for selenium. As appears from the present calculations of the pressure effect on $\chi$ for FeSe$_{0.5}$Te$_{0.5}$ and the earlier similar calculations for FeSe and FeTe, the large positive pressure effect on susceptibility in FeSe$_{1-x}$Te$_x$ compounds is determined by the dominating positive contribution caused by the strong sensitivity of paramagnetic susceptibility to internal structural parameter $Z$ and its change under pressure. It should be noted that the largest pressure effect on $\chi$ appears in FeTe compound, and that is a source of the observed its ferromagnetic state at high pressures.

Finally, the revealed here opposite trends in composition and pressure dependencies of superconducting transition temperature and magnetic susceptibility in the normal state indicate to antagonistic interrelation between superconductivity and magnetism in FeSe$_{1-x}$Te$_x$ chalcogenides. This tendency obviously has to be taken into account in further studies of possible role of magnetic excitations in the mechanism of superconductivity in Fe-based HTSCs.

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[1] M.D. Lumsden and A.D. Christianson, J. Phys.: Condens. Matter 22, 203203 (2010).
[2] C.W. Chu, B. Lorenz, Physica C 469, 385 (2009).
[3] J. Paglione and R.L. Greene, Nature Phys. 6, 645 (2010).
[4] J. Wen, G. Xu, G. Gu, J.M. Tranquada and R.J. Birgeneau, Rep. Prog. Phys. 74, 124503 (2011).
[5] L. Mazin, Nature 464, 183 (2010).
[6] P.J. Hirschfeld, M.M. Korshunov and I.I. Mazin, Rep. Prog. Phys. 74, 124508 (2011).
[7] Y. Kohama, Y. Kamihara, H. Kawai, T. Atake, and H. Hosono, Phys. Rev. B 78, 020512(R) (2008).
[8] M.V. Sadoskii, E.Z. Kuchinskii, I.A. Nekrasov, J. Magn. Magn. Mater. 324, 3481 (2012).
[9] Y. Mizuguchi and Y. Takano, J. Phys. Soc. Japan 79, 102001 (2010).
[10] D. Braithwaite, B. Salce, G. Lapertot, F. Bourdarot, C. Marin, D. Aoki, M. Hanfland, J. Phys.: Condens. Matter 21, 232202 (2009).
[11] S. Medvedev, T.M. McQueen, I.A. Troyan, T. Palasyuk, M.I. Eremets, R.J. Cava, S. Naghavi, F. Casper, V. Ksenofontov, G. Wortmann and C. Felser, Nature Mater. 8, 630 (2009).
[12] K. Horigane, N. Takeshita, C.-H. Lee, H. Hiraka, and K. Yamada, J. Phys. Soc. Japan 78, 063705 (2009).
[13] J. Pietosa, D.J. Gawryluk, R. Puzniak, A. Wisniewski, J. Fink-Finowicki, M. Kozlowski, and M. Berkowski, J. Phys.: Condens. Matter 24, 265712 (2012).
[14] H. Okada, H. Takahashi, Y. Mizuguchi, Y. Takano, and H. Takahashi, J. Phys. Soc. Japan 78, 083709 (2009).
[15] H. Takahashi, H. Okada, H. Takahashi, Y. Mizuguchi, and Y. Takano, J. Phys.: Conf. Series 200, 012196 (2010).
[16] B.C. Sales, A.S. Sefat, M.A. McGuire, R.Y. Jin, D. Manardus, and Y. Mozharivskyj, Phys. Rev. B 79, 094521 (2009).
[17] G.F. Chen, Z.G. Chen, J. Dong, W.Z. Hu, G. Li, X.D. Zhang, P. Zheng, J.L. Luo, and N.L. Wang, Phys. Rev. B 79, 140509 (2009).
[18] R. Viennois, E. Giannini, D. van der Marel, R. Černý, J. Solid State Chem. 183, 769 (2010).
[19] J. Yang, M. Matsu, M. Kawa, H. Ohta, C. Michioka, C. Dong, H. Wang, H. Yuan, M. Fang, and K. Yoshimura, J. Phys. Soc. Japan 79, 074704 (2010).
[20] T. Noji, T. Suzuki, H. Abe, T. Adachi, M. Kato, and Y. Koike, J. Phys. Soc. Japan 79, 084711 (2010).
[21] Y. Liu, R.K. Kremer and C.T. Lin, Supercond. Sci. Technol. 24, 035012 (2011).
[22] A.V. Fedorchenko, G.E. Grechnev, V.A. Desenko, A.S. Panfilov, S.L. Gatnchenko, V.V. Tsurkan, J. Deisenhofer, H.-A. Krug von Nidda, A. Loidl, D.A. Chariev, O.S. Volkova, A.N. Vasiliev, Low Temp. Phys. 37, 83 (2011).
[23] G.E. Grechnev, A.S. Panfilov, A.V. Fedorchenko, V.A. Desenko, S.L. Gatnchenko, V. Tsurkan, J. Deisenhofer, A. Loidl, D.A. Chariev, O.S. Volkova, A.N. Vasiliev, J. Magn. Magn. Mater. 324, 3460 (2012).
[24] Y.A. Ovchenkov, D.A. Chariev, E.S. Kozlyakova, O.S. Volkova, Physica C 489, 32 (2013).
[25] G.E. Grechnev, A.S. Panfilov, V.A. Desenko, A.V. Fedorchenko, S.L. Gatnchenko, D.A. Chariev, O.S. Volkova and A.N. Vasiliev, J. Phys.: Condens. Matter 25, 046004 (2013).
[26] A.V. Fedorchenko, G.E. Grechnev, V.A. Desenko, A.S. Panfilov, S.L. Gatnchenko, V. Tsurkan, J. Deisenhofer, A. Loidl, O.S. Volkova and A.N. Vasiliev, J. Phys.: Condens. Matter 23, 325701 (2011).
[27] D. Chariev, E. Osadchii, T. Kuznieceva, J.-Y. Lin, S. Kuznicheva, O. Volkova and A. Vasiliev, Cryst. Eng.
Comm. 15, 1989 (2013).

[28] M. Baran, V. Dyakonov, L. Gladczuk, G. Levchenko, S. Piechota, H. Szymczak, Physica C: Superconductivity 241, 383 (1995).

[29] L.D. Jennings and C.A. Swenson, Phys. Rev. 112, 31 (1958).

[30] J.M. Wills, M. Alouani, P. Andersson, A. Delin, O. Eriksson, A. Grechnev, Full-Potential Electronic Structure Method. Energy and Force Calculations with Density Functional and Dynamical Mean Field Theory. Springer Series in Solid-State Sciences, Springer Verlag, Berlin, Vol. 167, 200 p. (2010).

[31] G.E. Grechnev, Low Temp. Phys. 35, 638 (2009).

[32] U. von Barth and L. Hedin, J. Phys. C: Solid State Phys. 5, 1629 (1972).

[33] G. Tsoi, A.K. Stemshorn, Y.K. Vohra, P.M. Wu, F.C. Hsu, Y.L. Huang, M.K. Wu, K.W. Yeh, and S.T. Weir, J. Phys.: Condens. Matter 21, 232201 (2009).

[34] V. Tsurkan, J. Deisenhofer, A. Günther, Ch. Kant, M. Klemm, H.-A. Krug von Nidda, F. Schrettle, and A. Loidl, Eur. Phys. J. B 79, 289 (2011).

[35] P.S. Malavi, S. Karmakar, N.N. Patel, and S.M. Sharma, arXiv:1308.3367 [cond-mat.mtrl-sci] (2013).

[36] Y. Koshika, T. Usui, S. Adachi, T. Watanabe, K. Sakano, S. Simayi, and M. Yoshizawa, J. Phys. Soc. Japan 82, 023703 (2013).

[37] M. Bendele, A. Ichsanow, Yu. Pashkevich, L. Keller, Th. Strässle, A. Gusev, E. Pomjakushina, K. Conder, R. Khasanov, and H. Keller, Phys. Rev. B 85, 064517 (2012).

[38] Y. Mizuguchi, F. Tomioka, K. Deguchi, S. Tsuda, T. Yamaguchi, Y. Takano, Physics C 470, S353 (2010).

[39] T. Imai, K. Ahilan, F.L. Ning, T.M. McQueen, and R.J. Cava, Phys. Rev. Lett. 102, 177005 (2009).

[40] M. Bendele, A. Maisuradze, B. Rosésl, S.N. Gvasaliya, E. Pomjakushina, S. Weyeneth, K. Conder, H. Keller, and R. Khasanov, Phys. Rev. B 87, 060409 (2013).