Nature of extrinsic and intrinsic self-trapping of charge carriers in underdoped cuprate high-\(T_c\) superconductors

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Nature of extrinsic and intrinsic self-trapping (ST) of charge carriers in cuprates have been studied theoretically. The binding energies and radii of the extrinsic and intrinsic large polarons and bipolarons in cuprates are calculated variationally using the continuum model and adiabatic approximation. We have shown that the extrinsic and intrinsic three-dimensional (3D) large bipolarons exist in underdoped cuprates at \(\eta = \varepsilon_\infty/\varepsilon_0 < 0.127\) and \(\eta < 0.138\), respectively [where \(\varepsilon_\infty/\varepsilon_0\) is the optic (static) dielectric constant].

Key words: polaron, bipolaron, self-trapping, high-\(T_c\) superconductors

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1. Introduction

It is well known that electron-phonon coupling (EPC) is one of the most common and fundamental interactions in solids. In particular, a strong EPC in high-temperature superconducting cuprates (HTSC) was reported by a lot of experiments, which implies that EPC plays an important role in pairing. Accordingly, charge carriers in polar materials interacting with the phonon field can cause the formation of self-trapped polarons and bipolarons. The interest to the polarons and bipolarons is caused by their important role in explaining many characteristics of HTSC compounds (see [1,2] and references therein). The ground state of an undoped HTSC is an antiferromagnet (AF) whose doping by holes leads to the phase showing high temperatures of the superconducting (SC) transition. The extended \(t-t'\)-\(t''\)-\(J\)-Hubbard model, where a hole in a two-dimensional (2D) AF is simultaneously magnetic and a lattice polaron that moves through the lattice emitting and absorbing magnons and phonons, is used in many papers for calculation of spectral properties of underdoped HTSC compounds [3]. According to the band calculations and angle resolved-photoemission spectroscopy (ARPES) data [4,5], the electronic structure of the parent cuprate compounds is well described by a three-band Hubbard model and the oxygen valence band lies within the Mott-Hubbard gap. Furthermore, according to the combined study of the ARPES and quantum Monte Carlo simulations, there was proposed a novel polaronic metallic state in underdoped cuprates [6]. The variety of opinions range from a complete negation of the role of EPC in the physics of HTSC [6,7] to the statement that the binding energy of the polaron is an order of magnitude larger than any characteristic energy of the magnetic subsystem, and, therefore, exactly the magnetic system is irrelevant [8,9]. Both of the above radical statements have been criticized many times. On the other hand, the success of the extended \(t-J\)-Hubbard model in describing the ARPES [3,10] and optical spectra [11] does not make it possible to completely disregard the role of the magnetic subsystem. Correspondingly, there is an opinion that interactions with both magnetic and lattice subsystems are important [12]. Results indicating an important role of the EPC in HTSC compounds are presented, with emphasis on its implications for ARPES and optical conductivity [12].

Hole doping of the cuprates produces first quasi-free holes having the mass \(m_h\) in the oxygen valence band. The hole carriers are assumed to be within both a three-dimensional (3D) and a 2D deformable
medium, the last one being CuO$_2$ layers [13]. In reality, however, no systems can be purely 2D, and therefore, the layered cuprate compounds may be approximated as a 3D deformable medium. There is also a convincing experimental evidence that the consideration of cuprates as 3D systems may appear to be more appropriate (see [14–17]). The experimental results presented in [18] indeed confirm that the hole-doped system La$_{2−x}$Sr$_x$CuO$_4$ (LSCO) becomes less 2D in the strongly localized state. In polar materials, the hole carriers interacting both with lattice vibrations (i.e., acoustic and optical phonons) and with lattice defects (e.g., dopants or impurities), can easily be self-trapped near the defects and in a defect-free deformable lattice. Actually, the carrier localization in the cuprates is interrelated and the quantitative theory of this phenomenon is still lacking. Particularly, possible roles of large- and small-radius dopants, dopant-driven and carrier-driven inhomogeneities, carrier–defect–lattice and carrier–lattice interactions and other factors are very important for the localization of carriers in hole-doped cuprates. In the present paper we study the extrinsic (defect-assisted) and intrinsic (phonon-assisted) single particle and pair self-trapping (ST) of carriers using the continuum model of ionic crystal and adiabatic approximation. The possibility of the formation of localized extrinsic and intrinsic (bi)polaronic states as well as hydrogen-like impurity states in charge-transfer (CT) gap of the cuprates is examined and compared with the experimental data.

2. Calculation of the ground-state energy of the system of a defect-bound hole carrier in the polar crystal

Electron-phonon interaction (EPI) affects the electronic properties of semiconductors and polar crystals in various ways depending on the strength of the electron-phonon interaction. Among them, the polaron formation and the ST are its typical and important effects on the carriers. The relevant charge carriers in hole-doped cuprates are large polarons [13, 18] and the strong EPIs are responsible for the polaron formation and the ST are its typical and important effects on the carriers. The relevant charge carriers in hole-doped cuprates are large polarons [13, 18] and the strong EPIs are responsible for the formation of localized extrinsic and intrinsic (bi)polaronic states as well as hydrogen-like impurity states in charge-transfer (CT) gap of the cuprates is examined and compared with the experimental data.

We use a continuum model proposed in [20] and adiabatic approximation to calculate the ground-state energy of an interacting system of a defect (dopant)-bound hole carrier in the polar crystal. In the continuous approximation, the functionals of the total energies of the single-carrier and two-carrier systems can be written as follows:

$$E_1[\Psi(r)] = \frac{\hbar^2}{2m^*} \int [\nabla \Psi(r)]^2 \, d^3r - \frac{e^2}{2\epsilon} \int \frac{\Psi(r)\Psi^*(r')}{|r-r'|} \, d^3r \, d^3r'$$

$$- \frac{E_p^2}{2K} \int \Psi^4(r) \, d^3r - \frac{Z e^2}{\epsilon_0} \int \frac{\Psi^2(r)}{r} \, d^3r + \left( V_0 - \frac{E_d E_{D0}}{K} \right) \int \Psi^2(r) \delta(r) \, d^3r$$

(2.1)

and

$$E_2[\Psi(r_1, r_2)] = \frac{\hbar^2}{2m^*} \int [\nabla_1 \Psi(r_1, r_2)]^2 + [\nabla_2 \Psi(r_1, r_2)]^2 \, d^3r_1 \, d^3r_2 + \frac{e^2}{\epsilon_0} \int \frac{\Psi^2(r_1, r_2)\Psi^2(r_3, r_4)}{|r_1-r_2|} \, d^3r_1 \, d^3r_2 \, d^3r_3 \, d^3r_4$$

$$- \frac{2e^2}{\epsilon} \int \frac{\Psi^2(r_1, r_2)\Psi^2(r_3, r_4)}{|r_1-r_3|} \, d^3r_1 \, d^3r_2 \, d^3r_3 \, d^3r_4$$

$$- \frac{2E_p^2}{K} \int \Psi^2(r_1, r_2) \Psi^2(r_3, r_4) \, d^3r_1 \, d^3r_2 \, d^3r_3 \, d^3r_4 - \frac{2Z e^2}{\epsilon_0} \int \frac{\Psi^2(r_1, r_2)}{r_1} \, d^3r_1 \, d^3r_2$$

$$+ 2 \left( V_0 - \frac{E_d E_{D0}}{K} \right) \int \Psi^2(r_1, r_2) \delta(r_1) \, d^3r_1 \, d^3r_2, \quad (2.2)$$
where $m^*$ and $e$ are carrier effective mass and charge, respectively, $r_1$ and $r_2$ are the position vectors of the carriers, $E_d$ and $E_{dB}$ are the deformation potentials of the carrier and the defect, respectively, $\tilde{\varepsilon}$ is given by $\varepsilon^{-1} = \varepsilon_{\infty}^{-1} - \varepsilon_e^{-1} = (1 - \eta)/\varepsilon_{\infty}$, $K$ is an elastic constant, $V_0$ is the short-range defect potential, $Z$ is the charge state of the defect. A large ionicity of the cuprates $\eta = \varepsilon_{\infty}/\varepsilon_0 \ll 1$, enhances the polar EPI and the tendency to polaron formation.

In order to minimize the functionals (2.1) and (2.2) with respect to $\psi(r)$ and $\Psi(r_1, r_2)$ we can take the trial wave functions as follows:

$$
\psi(r) = N_1 \exp[-(\sigma r)] \tag{2.3}
$$

and

$$
\Psi(r_1, r_2) = N_2[1 + \gamma(\sigma r_{12})] \exp[-(\sigma r_1 + r_2)], \tag{2.4}
$$

where $N_1 = \sigma^{3/2}/\sqrt{\pi}$ and $N_2 = \sigma^3/\pi \sqrt{K_1(\gamma)}$ are the normalization factors, $\sigma = \beta/\alpha_0$, $K_1(\gamma) = 1 + \frac{\pi^2}{4} \gamma + 6\gamma^2$ is the correlation coefficient, $\beta$ and $\gamma$ are the variational parameters that characterize the localization degree of carriers and the correlation in their motions, respectively. $r_{12} = |r_1 - r_2|$ is the distance between the carriers, $\alpha_0$ is the lattice constant. Substituting equations (2.3) and (2.4) into equations (2.1) and (2.2), and performing the integrations in equations (2.1) and (2.2), we obtain the following functionals

$$
E_1[\beta] = A \left[ \beta^2 - g_s(1 + b_3) \beta^3 - g_1 \left( 1 - \eta + \frac{16}{5} Z \eta \right) \beta \right] \tag{2.5}
$$

and

$$
E_2[\beta, \gamma] = 2A \frac{K_2(\gamma)}{K_1(\gamma)} \left[ \beta^2 - \frac{8}{5} g_1 \left( 2(1 - \eta) \frac{K_1(\gamma)}{K_2(\gamma)} + Z \eta \frac{K_3(\gamma)}{K_2(\gamma)} \beta \right) \right. \\
- 16 b_2 \left. \left( \frac{K_6(\gamma)}{K_2(\gamma)} + \frac{b_3 K_4(\gamma)}{16 K_2(\gamma)} \right) \beta^3 \right], \tag{2.6}
$$

where $A = h^2/2m^* \alpha_0^2$, $g_s = E_d^2/16\pi K_0^2 A$ and $g_1 = 5e^2/16\varepsilon_{\infty} \alpha_0 A$ are dimensionless short-range and long-range carrier-phonon coupling parameters, $b_2 = 16(E_{dB}/E_d - K\varepsilon_0/E_0^2)$ is short-range carrier-defect-phonon coupling parameter, and the correlation coefficients are analytical functions of the variational parameter $\gamma$:

$$
K_2(\gamma) = 1 + \frac{25}{8} \gamma + 4\gamma^2,
$$

$$
K_3(\gamma) = \frac{5}{8} \frac{1087}{216} \gamma + \frac{38237}{2304} \gamma^2 + \frac{67639}{2592} \gamma^3 + \frac{4293}{256} \gamma^4,
$$

$$
K_4(\gamma) = \frac{5}{8} \frac{185}{216} \gamma + \frac{4199}{1728} \gamma^2 + \frac{8591}{2592} \gamma^3 + \frac{477}{256} \gamma^4,
$$

$$
K_5(\gamma) = 2 + \frac{15}{2} \gamma + 9\gamma^2,
$$

$$
K_6(\gamma) = \frac{1}{4} \frac{185}{216} \gamma + \frac{4199}{1728} \gamma^2 + \frac{8591}{2592} \gamma^3 + \frac{477}{256} \gamma^4,
$$

$$
K_7(\gamma) = 1 + 3\gamma + 3\gamma^2.
$$

Minimization of the functionals (2.5) and (2.6) over the variational parameters $\beta$ and $\gamma$ would give the ground state energies $E_1^{\text{min}}(\beta)$ and $E_2^{\text{min}}(\beta, \gamma)$ of single and pair self-trapped carriers. As a function of $\beta$ and $\gamma$, these functionals have also got the maxima $E_1^{\text{max}}(\beta)$ and $E_2^{\text{max}}(\beta, \gamma)$. The other parameters $A$, $g_s$, $g_1$, $b_3$ and $Z$ entering equations (2.5) and (2.6) play different roles in the formation of self-trapped states and determine the nature of the localized states of carriers in doped polar materials. The values of the parameters $A$, $g_s$ and $g_1$ can be obtained using the experimental values of the parameters $\alpha_0$, $m^*$, $K$, $E_0$, $\varepsilon_0$ and the Fermi energy $E_F$ of the undoped cuprates. The parameters $b_3$ and $Z$ characterize the formation of extrinsic (i.e., non-hydrogenic) self-trapped states of carriers, whereas the parameters $g_s$ and $g_1$ characterizing the strengths of short- and long-range carrier-phonon interactions are responsible for the formation of intrinsic self-trapped states.

Using the equations (2.5) and (2.6), we calculate the energies of different localized in-gap states in the CT gap of the cuprates. In order to determine the nature of these in-gap states and the quasi-free to localized state transition, we distinguish different physical situations in these systems. One can make
interesting analyses of real systems based on the sign and magnitude of $b_s$. In hole-doped cuprates, the situations might be quite different for different types of dopants. In particular, the signs of the deformation potential constants $E_d$ and $E_{dd}$ for holes and small-radius defects are always positive, while $E_{dd}$ for large-radius defects is negative [20]. However, at present no information is available for the magnitudes of the parameters $E_{dd}$ and $V_0$. Therefore, the parameter $b_s$ in equations (2.3) and (2.5) can be considered as the free parameter. We consider first the possibility of the formation of localized in-gap states at single and pair ST of carriers near the small-radius dopants (with $E_{dd} > 0$ or $b_s > 0$) in La-based cuprates. In this case, both short and long range parts of the defect potential in equations (2.3) and (2.5) are attractive, so that the substitution of small-radius cations (e.g., Ca$^{2+}$ and Nd$^{3+}$ ions) for La$^{3+}$ ions in La$_2$CuO$_4$ and for Sr$^{2+}$ ions in LSCO leads to a combined defect- and phonon-assisted ST of hole carriers with the formation of localized single-carrier and two-carrier impurity states, which are extrinsic polaronic and bipolaronic (the so-called U-pairing) states. At $Z \neq 0$ and $b_s > 0$, the minima of $E_3(\beta)$ and $E_2(\beta, \gamma)$ correspond to the ground-states of the extrinsic large polaron and bipolaron, respectively, measured with respect to the top of the oxygen valence band. The binding energies of such extrinsic large polaron and bipolaron (or negative U center) are defined as $E_{pol} = |E_{1}^{\text{min}}(\beta)|$ and $E_{bip} = |E_{2}^{\text{min}}(\beta, \gamma) - 2E_{1}^{\text{min}}(\beta)|$, respectively. In 3D systems, there is generally a potential barrier between the large- and small-radius self-trapped states. The two states of the extrinsic large polaron are separated by a potential barrier, with activation energy $E_1^{\text{sp}} = E_1^{\text{max}}(\beta) - E_1^{\text{min}}(\beta)$ needed for the transition from the large-radius localized state to the small-radius one. The potential barriers $E_{2}^{A} = E_{2}^{\text{max}}(\beta, \gamma) - E_{2}^{\text{min}}(\beta)$ exists between the large and small-radius extrinsic bipolaronic states.

We now calculate the basic parameters of the extrinsic large (bi)polarons in La-based cuprates. At low temperature, the La-based cuprates are orthorhombic with the lattice parameter $a_0 = 5.4$ Å. According to the spectroscopy data, the Fermi energy of the undoped cuprates is about $E_F = 7$ eV [21]. To determine the value of the short-range carrier-phonon coupling constant $g_s$, we can estimate the deformation potential $E_d = (2/3)E_F$ [22]. For the cuprates, typical values of other parameters are $m^* = m_e/13$, $\epsilon_{\infty} = 3-5$ [23, 24], $K = 1.4 \times 10^{12}$ dyne/cm$^2$ [25], and $Z = 1$. The calculated values of $E_{pol}$, $E_{bip}$, $E_1^{A}$ and $E_2^{A}$ for $b_s = 1$ and different values of $\epsilon_{\infty}$ and $\gamma$ are presented in Table 1. From Table 1, we can see that the potential barriers separating the large- and small-radius extrinsic (bi)polaronic states are rather high. These high potential barriers prevent the formation of small extrinsic (bi)polarons in 3D cuprates. The defect- and phonon-assisted ST of large polaron and large bipolaron in La-based cuprates are shown in figures 1 and 2 respectively.

### Table 1. The calculated parameters of the extrinsic large polarons and bipolarons (with correlation between the pairing carriers) in 3D cuprates at $Z = 1$, $b_s = 1$ and different values of $\epsilon_{\infty}$ and $\gamma$.

| $\eta$ | $\epsilon_{\infty} = 3.5$ | $\epsilon_{\infty} = 4$ | $\epsilon_{\infty} = 4.5$ |
| --- | --- | --- | --- |
| $E_{pol}$, eV | $E_{bip}$, eV | $E_1^{A}$, eV | $E_2^{A}$, eV | $E_{pol}$, eV | $E_{bip}$, eV | $E_1^{A}$, eV | $E_2^{A}$, eV | $E_{pol}$, eV | $E_{bip}$, eV | $E_1^{A}$, eV | $E_2^{A}$, eV |
| 0.00 | 0.1135 | 0.0610 | 5.3373 | 5.8611 | 0.0863 | 0.0456 | 5.6962 | 6.4752 | 0.0781 | 0.0345 | 5.9807 | 6.7938 |
| 0.02 | 0.1240 | 0.0631 | 5.2128 | 5.7469 | 0.0943 | 0.0391 | 5.5849 | 6.3881 | 0.0741 | 0.0303 | 5.8801 | 6.7990 |
| 0.04 | 0.1345 | 0.0654 | 5.0991 | 5.6895 | 0.1024 | 0.0322 | 5.4742 | 6.3012 | 0.0705 | 0.0261 | 5.7803 | 6.6024 |
| 0.06 | 0.1460 | 0.0687 | 4.9964 | 5.6345 | 0.1115 | 0.0266 | 5.3641 | 6.2152 | 0.0671 | 0.0228 | 5.6804 | 6.4395 |
| 0.08 | 0.1564 | 0.0721 | 4.8843 | 5.5798 | 0.1208 | 0.0216 | 5.2552 | 6.2192 | 0.0645 | 0.0185 | 5.5815 | 6.3495 |
| 0.10 | 0.1709 | 0.0759 | 4.7748 | 5.5361 | 0.1292 | 0.0166 | 5.1468 | 6.0438 | 0.1019 | 0.0058 | 5.4831 | 6.3625 |
| 0.12 | 0.1839 | 0.0801 | 4.6653 | 5.5069 | 0.1396 | 0.1376 | 5.0391 | 5.9588 | 0.1096 | 0.0057 | 5.3854 | 6.2779 |

Another interesting question is in what way large-radius dopants in cuprates affect the carrier-phonon system, especially near such defects. This opposite situation is realized in LSCO or La$_{2-x}$Ba$_x$CuO$_4$ (LBCO), where the radius of Sr$^{2+}$ ions is larger than that of La$^{3+}$ ions [26], so that for Sr$^{2+}$ ion $Z = 1$, $E_{dd} < 0$ or $b_s < 0$. In this case, the short-range part of the impurity potential in equation (2.1) is repulsive. Therefore, one can treat it like a hard core. The hole-lattice interactions near the large-radius dopants in LSCO and LBCO are suppressed by this repulsive defect potential and hole carriers are localized at a distance from the dopants (i.e., hole-carriers are loosely bound to dopants by long-range Coulomb attraction). From these considerations, it follows that the hole-lattice interaction near the large-radius dopants is weak and the localized impurity state may be of a hydrogen-like character described by a rigid lattice model [27]. Therefore, we can consider the hydrogen-like impurity centers having the
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Bohr radius \( a_H = 0.529 \varepsilon_0 (m_e/m^*) \) Å and the ionization energy \( E_{I}^{0} = e^2/2\varepsilon_0 a_H \) in lightly doped LSCO and LBCO. A similar argument can be made for other cuprates, such as \( \text{La}_2\text{CuO}_4 + \delta \), \( \text{YBa}_2\text{Cu}_3\text{O}_7 - \delta \) (YBCO) and \( \text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8 + \delta \) (Bi-2212). In these systems, the doping centers are excess oxygen atoms which capture the electrons from the oxygen conduction band and form the acceptor centers, i.e., \( \text{O}^- \) ions. One can assume that such extra \( \text{O}^- \) ions just like \( \text{Sr}^{2+} \) ions may have negative \( E_dD \).

Figure 1 shows the variation of the ratio \( E_{bU}/2E_{pl} \) with \( \eta \) for \( b_s = 0.5 \) and 2.0 for the stability region of the extrinsic large bipolaron in 3D cuprates. One can find out that such bipolarons exist as long as \( \eta \) is less than the critical value \( \eta_c \) = 0.127 and the ratio \( E_{bU}/2E_{pl} \) reaches up to 0.287 (at \( \varepsilon_\infty = 3 \) and \( \eta \to 0 \)). We have determined the stability region of the extrinsic large bipolaron in cuprates and found that such bipolarons exist as long as \( \eta \) is less than the critical value \( \eta_c \) = 0.127 and the ratio \( E_{bU}/2E_{pl} \) reaches up to 0.287 (at \( \varepsilon_\infty = 3 \) and \( \eta \to 0 \)). Figure 4 shows the variation of the ratio \( E_{bB}/2E_{p} \) with \( \eta \) for \( \varepsilon_\infty = 3, 4 \) and 5 for real large bipolarons in cuprates. One can see that in 3D cuprates, the intrinsic large bipolarons can exist at \( \eta = \eta_c = 0.138 \) and the ratio \( E_{bB}/2E_{p} \) reaches up to 0.27 (at \( \varepsilon_\infty = 3 \) and \( \eta \to 0 \)). While the inset in figure 4 shows the variation of the ratio \( E_{bB}/2E_{p} \) [calculated from equations (2.5) and (2.6) at \( g_s = 0 \) and \( Z = 0 \)] with \( \eta \) for large optical bipolaron. Of course, besides the ratio \( \eta = \varepsilon_\infty/e_0 \), the values of the Fröhlich electron-phonon coupling constant \( \alpha \) are very important for the formation of large optical bipolarons. The long-range coupling of carriers with optical phonons is much stronger than their short-range coupling with acoustic phonons. Therefore, the long-range Fröhlich-type EPI in polar materials...
Figure 3. (Color online) The ratio of the binding energy of the extrinsic large bipolaron to twice the extrinsic large polaron binding energy as a function of $\eta$ for two values of $b_s$ in 3D cuprates. The inset shows the dependence of the ratio of the binding energy of the extrinsic large bipolaron to twice that of the extrinsic large polaron on $\eta$ for two values of $\varepsilon_\infty$ in 3D cuprates.

have been studied extensively [28,29], although the short-range deformation potential type interaction is also important and leads to new effects. The dimensionless EPI Fröhlich coupling constant is defined as

$$\alpha = \frac{e^2(1-\eta)}{2\varepsilon_\infty \hbar \omega_{\text{LO}} (2m^* \omega_{\text{LO}} \hbar)^{1/2}},$$

(2.7)

where $\omega_{\text{LO}}$ is the frequency of the longitudinal-optical (LO) phonon in an ionic crystal.

In polar materials, the formation of optical bipolarons is favored by larger values of $\alpha$ and by smaller values of $\eta$ [2,29], i.e., the optical (Fröhlich or Pekar) bipolarons exist only if $\alpha$ is greater than a critical value $\alpha_c$ and when $\eta < \eta_c$. Such 3D bipolarons can exist above rather high critical values $\alpha_c$, e.g., $\alpha_c = 7.3$ as found by Adamowski [30] and $\alpha_c = 6.8$ found by Verbist, Peeters and Devreese [31]. Further, the value

Figure 4. (Color online) The ratio of the binding energy of the real large bipolaron to twice that of the real large polaron as a function of $\eta$ for different values of $\varepsilon_\infty$ in 3D cuprates. The inset illustrates the ratio of the binding energy of the optical large bipolaron to twice that of the optical large polaron on $\eta$ for $\varepsilon_\infty = 3.0$ in 3D cuprates.
of \( a_c \) corresponding to the onset of the strong coupling regime is found to be \( a_c = 5.8 \) \(^{[32]}\), below which the formation of large optical bipolaron in 3D systems is unlikely. At a given value of \( \eta \), the value of \( a_c \) depends on \( m^*, \varepsilon_{\infty} \) and \( \omega_{LO} \). The values of \( \hbar\omega_{LO} \) in high-\( T_c \) cuprates range from 0.03 to 0.05 eV \(^{[13,33]}\). Then, according to equation (2.7), the values of \( a \) corresponding to these high-\( T_c \) materials with \( m^* = m_e, \varepsilon_{\infty} = 3 \) and \( \eta = 0.02–0.10 \) are equal to \( a = 4.96–6.95 \). Thus, the conditions for the formation of large optical bipolarons are more favorable in the cuprates with \( \varepsilon_{\infty} = 3, \eta = 0.02–0.06 \) and \( \hbar\omega_{LO} = 0.03–0.04 \) eV, at which conditions \( a > a_c = 5.8 \) \(^{[32]}\) and \( a > a_c = 6.8 \) \(^{[31]}\) are well satisfied.

We note here that the largest values of \( \eta_{\infty} = 0.079–0.14 \) and \( E_{FB}/2E_p = 0.22–0.25 \) as found in the literature \(^{[8,36]}\) (see also \(^{[29]}\)) were obtained for the optical bipolaron and do not correspond to specific substances. At the same time, the critical values of \( \eta \) below which the optical bipolarons can exist in cuprates were small enough \( (\eta_{\infty} = 0.040–0.055) \), as estimated in \(^{[29]}\). Our results are quite impressive in the sense that both the real and the optical bipolaron in 3D cuprates, can really exist for relatively large values of \( \eta \) (figure 4) and the large bipolarons are formed with the binding energies reaching up to 27% (at \( \eta = 0 \)) of the large polaron energy. The distinctive feature of the cuprates is their very large ratio of static to high-frequency dielectric constants. This situation is favorable for carriers attracted to polarization well created by the other ones or to Coulomb centers (dopants) to form 3D intrinsic or extrinsic large (bi)polarons and the ratios \( s_{BI}/s_{Pl} \) and \( s_{BI}/s_{Pl} \) depend on several parameters. In particular, \( E_{BI}, E_p, E_{dB} \) and \( E_{FB} \) would rapidly increase with \( \varepsilon_{\infty} \) decreasing from 5 to 3, while the ratios \( E_{BI}/E_{Pl} \) and \( E_{BI}/E_{Pl} \) increase more slowly as \( \varepsilon_{\infty} \) decreases. Interestingly, \( E_p \) is an increasing function of \( \eta \) (table 1), while \( E_p \) is a decreasing function of \( \eta \) (table 2). Furthermore, we find that both \( E_{BI} \) and \( E_{Pl} \) decrease with increasing \( \eta \). We also find that the ratio \( E_{BI}/E_{Pl} \) appreciably increases with \( b_0 \) as shown in figure 3. The calculated values of the radii of extrinsic large (bi)polarons \( R_1, R_{BI} \) (at \( Z = 1 \) and \( b_0 = 1 \)) and intrinsic large (bi)polarons, \( R_p, R_B \) (at \( Z = 0 \) and \( b_0 = 0 \)) for different values of \( \varepsilon_{\infty} \) and \( \eta \) are presented in table 3.

### Table 2. Calculated parameters of the real large polaron and bipolarons (with correlation between the pairing carriers) in 3D cuprates at different values of \( \varepsilon_{\infty} \) and \( \eta \).

| \( \eta \) | \( \varepsilon_{\infty} = 2.5 \) | \( \varepsilon_{\infty} = 4 \) | \( \varepsilon_{\infty} = 4.5 \) |
|---|---|---|---|
| \( E_p, \text{eV} \) | \( E_{FB}, \text{eV} \) | \( E_{FB}, \text{eV} \) | \( E_{FB}, \text{eV} \) |
| \( b_0, \text{Å} \) | \( b_0, \text{Å} \) | \( b_0, \text{Å} \) | \( b_0, \text{Å} \) |
| 0.00 | 0.1107 | 0.0383 | 27.275 | 19.989 |
| 0.02 | 0.1063 | 0.0484 | 27.398 | 20.259 |
| 0.04 | 0.1019 | 0.0589 | 27.521 | 20.531 |
| 0.06 | 0.0977 | 0.0699 | 27.645 | 20.834 |
| 0.08 | 0.0935 | 0.0715 | 27.769 | 21.118 |
| 0.10 | 0.0895 | 0.0736 | 27.893 | 21.404 |
| 0.12 | 0.0853 | 0.0756 | 28.017 | 21.691 |
| 0.14 | 0.0816 | 0.0776 | 28.142 | 21.980 |

### Table 3. Calculated values of the radii of intrinsic and extrinsic large (bi)polarons \( R_p, R_1, R_B \) and \( R_{BI} \) in 3D cuprates at different values of \( \varepsilon_{\infty} \) and \( \eta \).

| \( \eta \) | \( \varepsilon_{\infty} = 2.5 \) | \( \varepsilon_{\infty} = 4 \) | \( \varepsilon_{\infty} = 4.5 \) |
|---|---|---|---|
| \( R_p, \text{Å} \) | \( R_1, \text{Å} \) | \( R_B, \text{Å} \) | \( R_{BI}, \text{Å} \) |
| \( R_p, \text{Å} \) | \( R_1, \text{Å} \) | \( R_B, \text{Å} \) | \( R_{BI}, \text{Å} \) |
| 0.00 | 8.6695 | 15.049 | 8.929 | 14.803 |
| 0.02 | 8.7917 | 15.577 | 9.261 | 15.507 |
| 0.04 | 8.9813 | 14.146 | 7.5886 | 11.777 |
| 0.06 | 9.1791 | 14.756 | 7.2478 | 11.482 |
| 0.08 | 9.3854 | 15.413 | 8.9409 | 11.199 |
| 0.10 | 9.6069 | 16.121 | 8.6652 | 10.926 |
| 0.12 | 9.8261 | 16.887 | 8.4084 | 10.662 |
| 0.14 | 10.062 | 17.718 | 8.1488 | 10.488 |

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3. Discussion

We now make some remarks about the characteristic (i.e., binding) energies of large bipolarons in the cuprates. The binding energies of extrinsic and intrinsic bipolarons strongly depend on $\epsilon_\infty$, $\varepsilon_0$ and $\eta$. The values of $\epsilon_\infty$ and $\varepsilon_0$ observed in high-$T_c$ cuprates are 3–5 [4, 37, 38] and 33–105 [4, 39], respectively, so that the values of $\eta$ range from 0.028 to 0.15. Using the values of $\epsilon_\infty = 3$ and $\eta = 0.03$, we find $E_{bU} = 0.07 \text{ eV}$ and $E_{bB} = 0.061 \text{ eV}$. If we take other experimental values of $\epsilon_\infty = 4$ and $\varepsilon_0 = 50$ for the cuprates (see [2, 29]), then we obtain $E_{bU} = 0.017 \text{ eV}$ and $E_{bB} = 0.016 \text{ eV}$ at $\eta = 0.08$. Further, at $\epsilon_\infty = 5$ and $\eta = 0.08$ we find $E_{bU} = 0.0097 \text{ eV}$ and $E_{bB} = 0.0102 \text{ eV}$. Thus, the extrinsic and intrinsic bipolarons can be experimentally found in high-$T_c$ cuprates in the energy ranges $\sim 0.01$–$0.07 \text{ eV}$ and $\sim 0.01$–$0.06 \text{ eV}$, respectively.

The binding energies of large polarons and bipolarons are manifested in the excitation spectra of the hole-doped cuprates as the temperature-independent low-energy gaps or pseudogaps, which are different from the high-energy CT gaps ($\Delta_{CT} = 1.5$–$2.0 \text{ eV}$ [13]) of the cuprates.

It is of interest to compare our results with experimental data on localized in-gap states (or bands) and energy gaps (which are responsible for the existence of insulating phase and are precursors to the pseudogaps observed in the metallic state) in hole-doped cuprates. The above extrinsic and intrinsic (bi)polaronic states as well as hydrogenic impurity states emerge in the CT gap of the cuprates. In the experiments, these localized states are displayed as the in-gap states. One can see that the value of $E_p = 0.13 \text{ eV}$ obtained at $\epsilon_\infty = 4$ and $\eta = 0.1$ (table 1) is consistent with experimental data for lightly doped La$_2$CuO$_{4+\delta}$ [13]. The in-gap impurity band observed in this system at 0.13 eV might be associated with the extrinsic large polarons. While the values of $E_p = 0.096$–$0.105 \text{ eV}$ (table 2) obtained at $\epsilon_\infty = 3.5$ and $\eta = 0.04$–$0.06$ agree reasonably well with the large pseudogap value $\sim 0.1 \text{ eV}$ observed in LSCO [40].

One of the important experimental observations is that in LSCO, the flatband [41], which is $\sim 0.12 \text{ eV}$ below the Fermi energy for $x = 0.05$, moves upwards monotonously with increasing $x$, but the flatband is lowered as $x$ decreases and loses its intensity in the insulating phase. Apparently, the flatband observed by ARPES in the lightly doped LSCO ($x = 0.05$) is the energy band of large polarons since the effective mass of carriers obtained from the analysis of the ARPES spectra is about $2.1 m_e$ [41]. The values of $R_p$ (table 3) are also in good agreement with the experimental values of the radii of polarons which vary from 6 to 10 Å in cuprates [13].

4. Conclusions

We have studied the possible mechanisms of carrier localization in inhomogeneous hole-doped cuprates. The quantitative theory of the impurity-assisted and phonon-assisted single particle and pair ST of hole carriers in 3D lightly doped cuprates is developed within the continuum model and adiabatic approximation. The possible mechanisms for carrier localization lead to the formation of extrinsic large (bi)polaronic states, the hydrogenic impurity states (i.e., impurities with loosely bound free carriers or large polarons) and intrinsic large (bi)polaronic states in the CT gap of the cuprates. We have variationally calculated the binding energies and radii of the extrinsic and intrinsic large polarons and bipolarons, taking into account the short- and long-range parts of the carrier-defect-phonon and carrier-phonon interactions. We have determined the stability region of the extrinsic large bipolaron in cuprates and found that such bipolarons exist as long as $\eta$ is less than the critical value $\eta_c = 0.127$ and the ratio $E_{bU}/2E_p$ (where $E_p$ and $E_{bB}$ are the binding energies of the extrinsic large polaron and bipolaron, respectively) reaches 0.287 (at $\epsilon_\infty = 3$ and $\eta \rightarrow 0$). We have obtained the conditions for the real large bipolaron stability and estimated the values of $E_{bB}/2E_p$ in 3D cuprates (where $E_p$ and $E_{bB}$ are the binding energies of the intrinsic large polaron and bipolaron, respectively).

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Природа зовнішніх і внутрішніх самозахопних носіїв струму в слабкозлегованих купратних високотемпературних напівпровідниках

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Теоретично вивчено природу зовнішніх і внутрішніх самозахопних носіїв струму в купратах. Енергії зв'язування і радіуси зовнішніх і внутрішніх великих поляронів і біполяронів в купратах обчислено варіаційно, використовуючи неперерву модель та адіабатичне наближення. Ми показали, що зовнішні і внутрішні тривимірні великі біполярони існують в слабкозлегованих купратах, відповідно, при \( \eta = \varepsilon_\infty / \varepsilon_0 < 0.127 \) і \( \eta < 0.138 \) [де \( \varepsilon_\infty \) і \( \varepsilon_0 \) — оптична (статична) діелектрична стала].

Ключові слова: полярон, біполярон, самозахоплення, високотемпературні надпровідники