CONSTRaining STEstellar FEEDBACK: SHOCK-IONiZED GAS iN NEARBY STARBURST GALAXIES

Sungryong Hong1,2, Daniela Calzetti1, John S. Gallagher III3, Crystal L. Martin4, Christopher J. Conselice5, and Anne Pellerin6

1 Department of Astronomy, University of Massachusetts, Amherst, MA 01003, USA
2 National Optical Astronomy Observatory, 950 North Cherry Avenue, Tucson, AZ 85719, USA
3 Department of Astronomy, University of Wisconsin–Madison, 475 North Charter Street, Madison, WI 53706, USA
4 Physics Department, University of California, Santa Barbara, CA 93106-9530, USA
5 University of Nottingham, School of Physics and Astronomy, Nottingham NG7 2RD, UK
6 Department of Physics and Astronomy, Texas A&M University, 4242 TAMU, College Station, TX 77843-4242, USA

Received 2013 January 14; accepted 2013 August 27; published 2013 October 17

ABSTRACT

We investigate the properties of feedback-driven shocks in eight nearby starburst galaxies using narrow-band imaging data from the Hubble Space Telescope. We identify the shock-ionized component via the line diagnostic diagram [O iii] (λ5007)/Hβ versus [S ii] (λλ6716, 6731) (or [N ii] (λλ6583))/Hα, applied to resolved regions 3–15 pc in size. We divide our sample into three sub-samples: sub-solar, solar, and super-solar, for consistent shock measurements. For the sub-solar sub-sample, we derive three scaling relations: (1) $L_{\text{shock}} \propto \Sigma_{\text{SFR, HL}}^{0.92}$, and (2) $L_{\text{shock}}/L_{\text{tot}} \propto (L_{\text{H}}/L_{\odot})^{-0.65}$, where $L_{\text{shock}}$ is the Hα luminosity from shock-ionized gas, $\Sigma_{\text{SFR, HL}}$, the star formation rate (SFR) per unit half-light area, $L_{\text{tot}}$ the total Hα luminosity, and $L_{\text{H}}/L_{\odot}$ the absolute H-band luminosity from the Two Micron All Sky Survey normalized to solar luminosity. The other two sub-samples do not have enough number statistics, but appear to follow the first scaling relation. The energy recovered indicates that the shocks from stellar feedback in our sample galaxies are fully radiative. If the scaling relations are applicable in general to stellar feedback, our results are similar to those by Hopkins et al. for galactic superwinds. This similarity should, however, be taken with caution at this point, as the underlying physics that enables the transition from radiative shocks to gas outflows in galaxies is still poorly understood.

Key words: galaxies: interactions – galaxies: ISM – galaxies: starburst – ISM: structure

Online-only material: color figures

1. INTRODUCTION

The energy and momentum deposited by star formation activity into the interstellar medium (ISM), a.k.a. stellar feedback, is a major, but still not fully characterized, mechanism that governs the formation and evolution of galaxies. The stellar winds and supernova (SN) explosions in star-forming regions provide energy and momentum to the surrounding ISM, changing its thermodynamic and kinetic properties, and sometimes driving galactic scale outflows. Theoretical works (Chevalier & Clegg 1985, hereafter CC85; Silich et al. 2003; Suchkov et al. 1996) suggest that the large SN rate in a starburst causes the SN remnants (SNRs) to merge together before a significant amount of energy is radiated away, and to transfer most of the energy outside the starburst volume. Such thermalized energy is a main power source for driving superwinds from starburst regions.

The hot gas and bipolar winds predicted by the CC85 model are, however, too hot to be observed. We mainly observe, at X-ray and optical wavelengths, the rapid cooling zones that are most likely associated with the entrainment and mass loading of cooler ISM or the boundaries where phases mix (Cecil et al. 2002; Martin et al. 2002; Strickland et al. 2004). These still provide important information: by observing the locations and intensities of the outer wind shocks, we can make local estimates of wind energy densities and, therefore, constrain the wind power from the starbursts.

Simulations and models of the propagation of feedback energy on galactic scale (De Young & Heckman 1994; MacLow & Ferrara 1999; Strickland & Stevens 2000; Hopkins et al. 2012) indicate that star formation rate (SFR) and galaxy mass determine the feedback ability to drive galactic outflows, though predictions point at galaxies that are about 10–100 times less massive than those in which superwinds are observed. Although additional complications may arise from the presence of active galactic nuclei (AGNs) in massive galaxies, these results demonstrate that our understanding of how the mechanical energy from star formation interacts with the surrounding gas and how efficient such energy is at driving galactic scale winds is still limited.

In this paper, we investigate the properties of feedback-driven shocks in eight nearby, AGN-free, starburst galaxies as a function of SFR and stellar mass: He 2–10, Holmberg II, NGC 1569, NGC 3077, NGC 4214, NGC 4449, NGC 5236 (M83), and NGC 5253. There are two main ionizing radiation fields in starburst galaxies: the radiation from stars and the cooling radiation from shocks. The gas ionized by radiative shocks (shock-ionized
gas) generally forms a thin and faint gas layer, while the gas ionized by stellar photons (photoionized gas) is the dominant ionized gas component in starburst galaxies and is morphologically diffuse. Because of this characteristic, the luminosity contrast between shock-ionized gas and photoionized gas is low in low-resolution ground-based observations, and the shock-ionized component is thus generally difficult to detect. The high angular resolution of the Hubble Space Telescope (HST), therefore, is necessary for separating shock-ionized gas from photoionized gas in external galaxies, not only from the line ratios but also in terms of morphology. HST images in optical narrow-band filters offer the opportunity to probe both shock-ionized and photoionized gas in our sample galaxies, over spatial scales in the 3–15 pc range, which is small enough to enable separation of the two ionized gas constituents. The stringent requirement for high angular resolution is such that the results presented in this paper have been mostly unexplored before.

2. GENERAL DATA ANALYSIS APPROACH

To identify the emission from shocks, we use the emission line diagnostic diagram: [O iii] (λ5007)/Hβ versus [S ii] (λλ6716, 6731) (or [N ii] (λ6583))/Hα. The original use of this diagnostic was to discriminate starburst from AGN activity in galaxies, since the line ratios are sensitive to the hardness of the ionizing radiation field (Baldwin et al. 1981; Veilleux & Osterbrock 1987; Kewley et al. 2001, hereafter K01; Kauffmann et al. 2003). For this goal, galaxy-averaged line ratios are plotted on the diagnostic diagram, each point representing one galaxy.

As a different application of the same diagram, Calzetti et al. (2004, hereafter C04) plotted the line ratios of individual regions (bins of 5–10 pc size) from the spatially resolved images of four nearby starburst galaxies, to separate the shock-ionized component from the photoionized gas in the galaxies’ ISM. They adopted the “maximum starburst line” (MSL) from K01 as a shock separation criterion; the MSL is a theoretical limit, where the region to the upper right of the line cannot be explained by photoionization alone. C04 find that the estimated Hα luminosity from shock-ionized gas is a few percent of the total Hα luminosity. Their calculations indicate that the mechanical energy from the central starburst is enough to power the radiative shocks and a significant amount of the mechanical energy is radiated away. Hong et al. (2011, hereafter H11) adopt a larger range of properties for the photo- and shock-ionization models and apply various shock separation criteria to NGC 5236 (M83). The estimated shock Hα luminosity can vary from a couple of percent (from the “MSL”; the most conservative criterion) to 30% (from the most generous criterion [S ii] (λλ6716, 6731)/Hα > −0.5) of the total Hα luminosity. An intermediate estimate of the Hα luminosity in shocks places it at about 15% of the total, which implies that virtually all of the mechanical energy from the starburst is radiated away.

Building upon the previous studies of C04 and H11, we analyze in this paper a larger sample of eight galaxies and map out the relations between the feedback-driven shocks and global galactic parameters: SFR (≈mechanical energy injection rate) and host galaxy stellar mass (≈gravitational potential depth). With this, we attempt to specify relations that can inform models of galaxy formation and evolution. In Section 3, we describe the sample and data reduction processes. In Section 4, we describe the diagnostic diagrams and their related ionization models. Then, we present our main results. In Section 5, we discuss and summarize our results.

3. DATA DESCRIPTION

3.1. Sample Description

In order to secure a sample of actively star-forming galaxies spanning a range in stellar mass and SFR, the following criteria were applied: (1) distance < 12 Mpc (to exploit the HST angular resolution, 0′.2 = 12 pc at 12 Mpc, matched to the observed width of shock fronts in other galaxies; see Section 2 of C04); (2) recession velocity <950 km s⁻¹ (to get the emission lines inside the available narrow-band filters); and (3) centrally concentrated star formation/starburst.

Observations were performed with the HST through a number of programs (GO-9144, 10522, 11146, 11360) that gathered narrow-band images in F502N, F656N (or F658N, F657N), F487N, and F673N, centered on the relevant emission lines, plus two broad-band images in F547M (or F555W, or F550M) and F814W for stellar continuum subtraction. A total of eight starburst galaxies were observed via those GO programs, and a recently approved observing program will secure an additional two (Mrk 178 and NGC 4861, with program GO-12497). As a general operational approach, we produce emission line images for Hα, Hβ, [O iii] (λ5007), and [S ii] (λλ6716, 6731) (or [N ii] (λ6583)) from the narrow images by subtracting stellar continuum using broad-band images.

Figure 1 shows the distribution of our eight galaxies in the parameter space of absolute H-band magnitude from the Two Micron All Sky Survey (2MASS) (a proxy for stellar mass) versus the SFR within the central starburst region. Table 1 summarizes the general information about the sample galaxies and the instruments and filters used in the observations. We categorize the sample into three sub-groups according to metallicity: sub-solar (solid green circles in Figure 1; galaxies with metallicity around Z = 0.4 Z⊙), solar (open blue triangles; galaxies with metallicity around Z = 1.0 Z⊙), and...
super-solar (open purple square; galaxies with metallicity near $Z = 2.0 Z_\odot$). This grouping is guided by the metallicity dependence of the line ratios, which thus affects the numerical definition of shocked gas. More details on this effect are given in Section 4.1.

### 3.2. Data Reduction

Our sample was observed using the three HST instruments WFPC2, ACS, and WFC3. Table 2 summarizes the narrow-band imaging observations: filter, instrument, target emission line, program ID, exposure time, and sensitivity. For ACS and WFC3, we use the standard pipeline image products, which include processing through the MultiDrizzle software (Fruchter et al. 2009). For WFPC2, the STScI pipeline processing includes only basic steps such as flat-fielding and bias subtraction, for which we use the best reference files available in each observing period (Gonzaga & Biretta 2010). Our post-pipeline steps thus include removal of warm, hot pixels and cosmic rays, and registration and co-addition of the multiple images in each band. We use the IRAF (Image Reduction and Analysis Facility) task WARMPIX to remove hot pixels and the task CRREJ to remove cosmic rays and combine the dithered images (see C04 for more details about the WFPC2 post-pipeline processing).

Then, we perform the four steps below:

1. Photometric calibration.
2. Stellar continuum subtraction from narrow-band image.
3. Decontamination of [N ii] ($\lambda\lambda 6548, 6583$) from narrow-band image for H$\alpha$ if necessary.
4. Dust extinction correction.

We use the keyword PHOTFLAM for the photometric calibration of each galaxy in each band. The filters F487N and F502N only include a single emission line, H$\beta$ and [O iii] ($\lambda 5007$), respectively. This makes the photometric calibration simpler than the redder filters. For the [S ii] ($\lambda\lambda 6716, 6731$) calibration, we adopt the doublet ratio, [S ii] ($6716/6731$) $\approx 1.2$. This density indicator varies from 1.0 to 1.4 for most ISM conditions, and, within this range, the calibrated fluxes change by less than 10%. The uncertainty from continuum subtraction is generally larger than this uncertainty. For Holmberg II and NGC 4449, we observe the [N ii] ($\lambda 6583$) line using the F660N filter, instead of the [S ii] ($\lambda\lambda 6716, 6731$) line. The [N ii] ($\lambda 6583$) decontamination of the H$\alpha$ filter, thus, is straightforward for these two galaxies.

The second and third steps produce most of the uncertainty when calculating the line ratios for the diagnostic diagram. The fourth step, the dust extinction correction, typically does not affect the line ratios in the diagnostic diagram in a significant manner, due to the proximity in wavelength of the lines in each ratio, under the assumption that the line emission is coming from the same spatial region. However, the dust corrections for each line luminosity can be large. The following subsections describe in detail the impact of each of the last three steps, and how we deal with them.

#### 3.2.1. Stellar Continuum Subtraction

For each narrow-band image, we approximate the stellar continuum baseline near the target emission line using broad-band images, straddling, when possible, a line with two adjacent broad-band filters. Specifically, we use F547M (or F555W, or F550M) as reference stellar continuum for the H$\beta$ and [O iii] ($\lambda 5007$) and an interpolated continuum for H$\alpha$ and [S ii] ($\lambda\lambda 6716, 6731$) (or [N ii] ($\lambda 6583$)) using F547M and F814W. For the F555W filter, we remove the
self-contamination due to Hβ and [O III] (λ5007) line emission within the broad-filter bandpass using the iterative method described in H11. To find the optimal subtraction, we apply the skewness transition method to all of our narrow-band images (S. Hong et al. 2013, in preparation). This method is based on a feature that appears in the skewness at the transition between over- and under-subtraction of the stellar continuum from narrow-band images. The stellar continuum subtraction step is the one that produces the largest uncertainty in our final line emission images, due to the intrinsic limitations of the method, which cannot account, e.g., for color differences among stellar populations across the filter bandpasses.

### 3.2.2. [N II] Correction in the Hα Filter

Except for Holmberg II and NGC 4449, for which the [N II] line was directly observed, assumptions have to be made in order to remove the [N II] contamination from the narrow-band images targeting the Hα emission. We use two methods to deal with this problem. The first is to use a line ratio [N II]/Hα obtained from spectroscopy (e.g., from the literature). Because we use a single value for the whole image of each galaxy, spatial variation of the [N II]/Hα ratio cannot be considered in this method. The second is to use the relation [N II] ∝ [S II], which is less dependent on variations in the metal abundance and/or UV radiation. With this method, we still assume a constant factor for the entire image. Therefore, the [N II] correction is also a limitation for the photometric calibration of the Hα images in our sample, although the impact of [N II] variations is expected to be significant (>10%) only for M83, which has a large [N II]/Hα ratio. Table 1 summarizes the [N II] correction method applied to each galaxy.

#### 3.2.3. Dust Extinction Correction

We produce extinction maps using the Hα/Hβ ratios and the standard extinction equation

\[
E(B - V) = \frac{\log[H\alpha/H\beta]_a - \log[H\alpha/H\beta]}{0.4(k_{H\beta} - k_{H\alpha})},
\]

From Cardelli et al. (1989), we choose a normalization, \( R_V = 3.1 \), and use the Milky Way extinction curve; \( k_{H\alpha} = 3.609 \), \( k_{5890} = 3.473 \), \( k_{5563} = 2.535 \), \( k_{5653} = 2.525 \), and \( k_{6725} = 2.458 \), following the recipe of Calzetti (2001). This approach enables us to include both internal and foreground (from our own Milky Way) extinction simultaneously. The dust correction for the line ratios on the diagnostic diagram can be written as

\[
\log \left( \frac{[S\ II]}{H\alpha} \right)_i = \log \left( \frac{[S\ II]}{H\alpha} \right)_a - 0.0310E(B - V),
\]

\[
\log \left( \frac{[N\ II]}{H\alpha} \right)_i = \log \left( \frac{[N\ II]}{H\alpha} \right)_a - 0.0040E(B - V),
\]

\[
\log \left( \frac{[O\ III]}{H\beta} \right)_i = \log \left( \frac{[O\ III]}{H\beta} \right)_a - 0.0542E(B - V),
\]

where the “i” subscripts represent the intrinsic line ratio and the “a” subscripts represent the observed attenuated line ratio. The coefficients preceding the color excess \( E(B - V) \) are sufficiently small that dust corrections are in general a minor effect on the diagnostic diagram. They are needed only for extreme cases (e.g., for the heavy foreground extinction affecting NGC 1569).

### Table 2

**Summary of the Narrow-band Observations**

| Galaxy | \( z \) | Filter | Line | Continuum | Program ID | EXPTIME (s) | 1σ limit (erg s\(^{-1}\) cm\(^{-2}\)) |
|--------|--------|--------|------|-----------|------------|-------------|----------------|
| He 2–10 | 0.002912 | F658N(WFPC2) | Hα + [N II] | F547M F814W | 11146 | 2000 | 6.3 \times 10^{-18} |
| F487N(WFPC2) | Hβ | F547M | 11146 | 4400 | 1.0 \times 10^{-17} |
| F502N(WFPC2) | [O III] | F547M | 11146 | 5200 | 8.0 \times 10^{-18} |
| F673N(WFPC2) | [S II] | F547M F814W | 11146 | 5200 | 4.7 \times 10^{-18} |
| NGC 1569 | −0.000347 | F656N(WFPC2) | Hα + [N II] | F547M F814W; 8133 | 8133* | 1600 | 1.2 \times 10^{-17} |
| F487N(WFPC2) | Hβ | F547M | 8133* | 2400 | 1.7 \times 10^{-17} |
| F502N(WFPC2) | [O III] | F547M | 8133* | 1500 | 2.1 \times 10^{-17} |
| F673N(WFPC2) | [S II] | F547M F814W | 8133* | 3000 | 1.3 \times 10^{-17} |
| NGC 4449 | 0.000690 | F658N(ACS) | Hα + [N II] | F550M F814W | 10585* | 1539 | 3.8 \times 10^{-18} |
| F660N(ACS) | [N II] + Hα | F550M F814W | 10522 | 1860 | 1.3 \times 10^{-17} |
| F487N(WFPC2) | Hβ | F550M | 10522 | 2100 | 8.0 \times 10^{-18} |
| F502N(ACS) | [O III] | F550M | 10522 | 1284 | 1.0 \times 10^{-17} |
| Ho-2 | 0.000474 | F658N(ACS) | Hα + [N II] | F550M F814W | 10522 | 1680 | 1.0 \times 10^{-17} |
| F660N(ACS) | [N II] + Hα | F550M F814W | 10522 | 1686 | 6.4 \times 10^{-18} |
| F487N(WFPC2) | Hβ | F550M | 10522 | 5400 | 6.6 \times 10^{-18} |
| F502N(ACS) | [O III] | F550M | 10522 | 1650 | 7.2 \times 10^{-18} |
| NGC 5236 | 0.001711 | F657N(WFPC3) | Hα + [N II] | F555W F814W | 11360 | 1484 | 6.1 \times 10^{-18} |
| F487N(WFPC3) | Hβ | F555W | 11360 | 2700 | 4.3 \times 10^{-18} |
| F502N(WFPC3) | [O III] | F555W | 11360 | 2484 | 4.6 \times 10^{-18} |
| F673N(WFPC3) | [S II] | F555W F814W | 11360 | 1850 | 5.8 \times 10^{-18} |
| NGC 4214 | 0.000970 | F657N(WFPC3) | Hα + [N II] | F555W F814W | 11360 | 1592 | 5.3 \times 10^{-18} |
| F487N(WFPC3) | Hβ | F555W | 11360 | 1760 | 5.8 \times 10^{-18} |
| F502N(WFPC3) | [O III] | F555W | 11360 | 1470 | 6.2 \times 10^{-18} |
| F673N(WFPC3) | [S II] | F555W F814W | 11360 | 2940 | 3.5 \times 10^{-18} |

**Notes.**

1. NGC 3077 and NGC 5253 are absent here because we adopt the results for the galaxies from Calzetti et al. (2004; GO–9144). The starred program IDs are open access, not related to our programs.
Hβ maps at low surface brightness levels can produce artificially high values of the color excess $E(B - V)$, as will be discussed below.

The statistical distributions of astronomical data are generally Poissonian or Gaussian. The ratio of two random variables with Poissonian or Gaussian uncertainty distributions can produce a skewed distribution of the ratios. If we assume two normal (Gaussian) distributions, expressed as $X = N(\mu_X, \sigma_X^2)$ and $Y = N(\mu_Y, \sigma_Y^2)$, where $\mu$ represents the mean and $\sigma$ represents the standard deviation, their ratio distribution, $Z = X/Y$, follows the equation below (Hinkley 1969):

$$
p_Z(z) = \frac{b(z)c(z)}{a^2(z)} \frac{1}{\sqrt{2\pi}\sigma_x\sigma_y} \left[2\Phi\left(\frac{z-\mu}{\sigma}\right) - 1\right] + \frac{1}{a^2(z)\pi\sigma_x\sigma_y} e^{-\frac{1}{2}\left(\frac{z-\mu}{\sigma}\right)^2},
$$

where

$$
a(z) = \sqrt{\frac{z^2}{\sigma_x^2} + \frac{1}{\sigma_y^2}},
$$

$$
b(z) = \frac{\mu_x}{\sigma_x^2} + \frac{\mu_y}{\sigma_y^2},
$$

$$
c(z) = e^{-\frac{z^2}{2\sigma_x^2}} - 1 - \frac{z^2}{2\sigma_x^2},
$$

$$
\Phi(z) = \int_{-\infty}^{z} \frac{1}{\sqrt{2\pi}} e^{-u^2/2} du.
$$

Since Hα and Hβ emissions are strong recombination lines, their intrinsic flux ratio is quite robust in most physical conditions: $H\alpha/H\beta \approx 2.87$. Because of this property, Equation (1) is a standard approach to measure the amount of extinction from the observed line ratios. However, Equation (1) is only valid when our observational sensitivity is infinitely high; in other words, the distribution of observed Hα/Hβ ratio must follow a delta-function, $\delta(z - 2.87)$, when there is no dust extinction. Indeed, the ratio distribution (Equation (2)), for sufficiently large $\mu_X$ and $\mu_Y$ with Poissonian variances ($\sigma_X^2 \sim \mu_X$), shows a delta-function-like distribution (see the probability distribution for $\mu_{H\beta} = 10^3$ in the top left panel in Figure 2). This means that for a bright region the dust extinction correction using Equation (1) is reliable. But, for a faint region, the Hα/Hβ ratios are heavily affected by statistical measurement errors.

The important point is that, even though there is no dust extinction, the Hα/Hβ line ratio can be different from the intrinsic value, 2.87, due to stochastic errors and shot noise. This is significant especially for diffuse gas (which is generally faint, implying large stochastic errors). The bottom left panel of Figure 2 shows the ratio distributions when the variance is dominated by background (BG) or instrumental noise (dark currents and readout noises; INST), while the top left panel shows the ratio distribution when the variance is dominated by the signal counts and the noise is Poissonian. Even for the 7σ detections for both cases, the distributions of line ratio are broad.

The right panel shows the Hβ flux versus the observed Hα/Hβ ratios for NGC 4214. Each point is one bin in the images as indicated in Table 1. NGC 4214 is one of the least dust-extincted galaxies in our sample, as verified from the galaxy-wide Hα/Hβ ratio, thus a perfect test case for the present discussion. As can be seen from the right panels of Figure 2, while the Hα/Hβ ratios span a small range at high Hβ fluxes in the pixel-by-pixel7 distribution (the filled histogram in the bottom right panel), the spread of Hα/Hβ ratios increases toward low Hβ flux values (the unfilled histogram), when we separate the pixels using our selected threshold value of 80 ($\approx 14\sigma$) for Hβ flux in the given scale unit. The threshold value of 80 is high enough, above which the stochastic broadening effect is small. While some of the scatter at low flux values may be due to intrinsic variations in the extinction values, we cannot exclude that a portion may be due to the statistical fluctuations presented in this section.

To show a possible bias due to the statistical fluctuations described above, we present the case of NGC 4449 in Figure 3; this galaxy has a non-negligible amount of internal dust extinction. The left panels show the probability functions for Hα/Hβ = 3.5 corresponding to $E(B - V) = 0.2$, which is the centroid of the filled histogram of the bottom right panel. In the left panels, the overall distributions look similar to the ones in Figure 2, but they become broader. The right panels show the observed ratios. Due to the shallower depth of the observations for NGC 4449, we set the Hβ threshold line at 150 ($\approx 16\sigma$) in the given scale unit. Because of the internal dust extinction, the observed ratios come from the convolution of the dust extinction distribution with its statistical broadening. For the brighter pixels above the threshold (the filled histogram in the bottom right panel), the “true” mean value is recovered, despite the broadening due to the statistical fluctuations. For the pixels below the threshold, the statistical fluctuations dominate the observed line ratios, and the peak of the distribution is below the “true” mean value. It is inevitable, for pixel-by-pixel correction, that we have negative extinctions when Hα/Hβ < 2.87. We set $E(B - V)$ to zero for those pixels. For galaxy-averaged measurements, the ratios can be averaged out by summing all of the pixels. However, for pixel-by-pixel correction, the step of imposing $E(B - V) = 0$ for otherwise negative extinction values causes overcorrection for dust extinction effects.

To summarize, the issue described in this section is present for pixel-by-pixel analyses of low surface brightness (high statistical uncertainty) regions; it is a minor effect in high surface brightness regions or in galaxy-averaged quantities. In light of the above, as a general rule, we will use the line ratios in the diagnostic diagrams without extinction corrections. However, such corrections will be important for line fluxes, and they will be applied. We will need to keep in mind, however, the caveats discussed in this section.

3.2.4. Error Bars and Most Uncertain Ratios (MURs) on Diagnostic Diagram

We impose a minimum threshold value of 5σ to our emission line images, in order to construct line ratios (see Section 4.4.2 for a discussion on changing the threshold). This implies that the uncertainty on the emission line ratios will be the largest when both emission lines are detected at the threshold value; any other line ratio will have lower uncertainty values, but will also come from brighter line value(s). The ratio involving two emission lines detected at the threshold value will thus determine a pivot ratio, with two properties: (1) the deepest detection (i.e., lowest flux value) and (2) the largest error bar. Hereafter, we call this the most uncertain ratio (MUR).

7 More correctly, bin-by-bin. To emphasize that our analysis is based on pixel-size scale, not on galactic scale, we use the term “pixel” for the bins used in this paper.
Figure 2. Probability density functions of Hα/Hβ ratio for zero extinction, i.e., the intrinsic Hα/Hβ = 2.87, from Equation (2) for various variances (left panels). For a bright source, the noise is dominated by the Poisson fluctuations on the source counts. The top left panel represents this source-dominated case, while the bottom left panel is for a background and instrument (BG+INST) noise-dominated case. The observed Hα/Hβ ratio vs. the Hβ flux for NGC 4214 (top right panel) and its number count histograms (bottom right panel). The vertical lines represent the positions of Hα/Hβ = 2.87. The horizontal dashed line on the top right panel shows the threshold line where the Hβ flux is equal to 80 in the given scale unit. Above the threshold line (the filled histogram in the bottom right panel), the signal is relatively strong; hence, the broadening caused by statistical measurement error is relatively small as shown in the top left panel (the probability function for μHβ = 10⁵). Below the threshold (the unfilled histogram in the bottom right panel), the statistical broadening becomes larger. Even though we consider other factors affecting the observed ratios, such as real dust extinction and residuals from stellar continuum subtraction, the statistical broadening shows a significant impact on the observed ratios. This statistical bias can propagate during dust extinction correction.

We locate the position of the MUR in each diagnostic diagram, along both axes. We call the intersection of the two MUR values the “MUR spot.” By adjusting the detection threshold or the observational depth, we can move the MUR spot. For example, if all images are observed with equal depth (in flux), the MUR spot is positioned at (0, 0) in log scale. If, conversely, [S ii] (λλ6716, 6731) and [O iii] (λ5007) are observed with 10 times the depth of the Hα and Hβ lines, respectively, the MUR spot moves to (−1, −1). This adjustment can be useful if we have a specific target ratio: for example, for shocks log([S ii] (λλ6716, 6731)/Hα) > 0.0. Figures 5 and 6 show the MUR spots with their error bars in our diagnostic diagrams; the error bars are FWHMs of ratio distributions. The thin (cyan) error bars are for 5σ detections, while the thick (blue) error bars are for 10σ detections, hence the latter are smaller than the former. Any other line ratio that is not on the MUR spot will have smaller error bars than those at the MUR spot.

4. RESULTS

4.1. Diagnostic Diagram and Metallicity Dependence

In order to establish whether the line emission from each of the subarcsecond bins in our galaxies is dominated by photoionized gas or shock-ionized gas, we need to compare the observed line ratios against models. For this purpose, we adopt the theoretical grids of K01 for photoionization and Allen et al. (2008, hereafter A08) for radiative shocks. The photoionization grids from K01 are based on the stellar population synthesis model STARBURST99 and the gas ionization code MAPPINGS III (Binette et al. 1985; Sutherland & Dopita 1993; Leitherer et al. 1999). The spectral energy distributions (SEDs) from STARBURST99 provide the ionizing fluxes, and the MAPPINGS III code calculates the ionization state for the atomic species and the fluxes of the emission lines. We derive line ratios from those model fluxes for a range of ionization parameter values (q, defined as the ratio of mean ionizing photon density to mean atom density in K01, ranges from 5.0 × 10⁶ to 3.0 × 10⁸) and for selected values of the metallicity and density of the gas. Figure 4 shows the photoionization tracks for a constant star formation history, Geneva stellar evolution tracks (Schaller et al. 1992), and Lejeune stellar atmosphere models (Lejeune et al. 1997). We only need these tracks to provide a consistent way to separate photoionized from shock-ionized gas, and not for quantitative analysis. Thus, we adopt the conservative track termed “MSL” (MSL in the legend of Figure 4) of K01 as our separating criterion between the two gas components. Above and to the right of this track, line ratios cannot be explained by photons from synthesized stellar populations.
In our case, we consider this non-stellar ionization to be shocks generated by stellar mechanical feedback.

The shock-ionization models of A08 calculate the ionizing radiation field from hot radiative shock layers dominated by free–free emission and use the MAPPINGS III code for the gas ionization state and the intensity of the emission lines. For radiative shocks, the emission comes from two components: the shock layer (post-shock component) and the precursor (pre-shock component). The shock layer is the cooling zone of the radiative shock, and the precursor is the ionized region by upstreaming photons from the cooling zone. Since the main radiation process in the shock layer is free–free emission, the ionizing radiation field from shocks is mainly determined by the shock Mach number (i.e., shock velocity), the pre-shock gas density, and the intensity of the ISM magnetic field (see A08 for more details). The ISM magnetic field affects the post-shock gas density; higher ISM magnetic fields result in lower post-shock gas densities that affect the ionization parameter of the post-shock gas component. For a given metallicity and pre-shock gas density, the line ratios are thus determined mainly by the shock velocity and, as a second main parameter, by the magnetic field. This is shown in Figure 4, for shock velocities from 200 km s$^{-1}$ to 500 km s$^{-1}$, with a minimum ISM magnetic field of $10^{-4}$ μG or 0.5 μG. The side branches from the selected shock velocities, 200, 250, ..., 500 km s$^{-1}$, show the effect of changing the magnetic field strength from $10^{-4}$ μG (or 0.5 μG) to 10 μG. This discussion assumes that the ambient photoionization rates are small. If that is not the case and the cooling zone is photoionized or if projection effects are important, some fraction of the shocked gas may be missed because of dilution.

Given the emission lines available for our sample galaxies, we use two diagnostics, [O iii] (λ5007)/Hβ versus [S ii] (λλ6716, 6731)/Hα (hereafter [S ii] diagnostics) and [O iii] (λ5007)/Hβ versus [N ii] (λ6583)/Hα (hereafter [N ii] diagnostics), as summarized in Table 1. Since the [O iii] (λ5007)/Hβ and [N ii] (λ6583)/Hα ratios strongly depend on metallicity (Kewley & Dopita 2002), the shock identification through the MSL is inevitably affected by metallicity. Figure 4 shows the theoretical tracks for the [S ii] diagnostics at four different metallicities, 0.2 $Z_{\odot}$, 0.4 $Z_{\odot}$, 1.0 $Z_{\odot}$, and 2.0 $Z_{\odot}$. The grids show the metallicity dependence of the [O iii] (λ5007)/Hβ ratio, especially conspicuous for $Z \geq 1.0 Z_{\odot}$. The shock tracks also show different patterns for each metallicity. This is easily seen by choosing a constant shock velocity, e.g., $v = 250$ km s$^{-1}$, which is a typical velocity for the observed shocks (see below), and following the variations of this track for changing metallicity. The most noticeable effect in Figure 4 is a decrease of the expected [O iii] (λ5007)/Hβ shock line ratio for increasing metallicity. The top panel in Figure 6 shows the theoretical tracks for the [N ii] diagnostic, which also presents a clear metallicity dependence of the [N ii] (λ6583)/Hα ratios.

Figure 5 shows the six galaxies for which the [S ii] diagnostic is used. From the distribution of pixels on the diagnostic
diagram, we can group the galaxies as (NGC 1569, NGC 5253, NGC 4214), (He 2–10, NGC 3077), and (NGC 5236). This grouping corresponds to similar oxygen abundance within each group (Table 1). An important consideration is that the pixel distribution of each galaxy roughly follows the expected trend at the galaxy’s metallicity value, as shown by the overplotted theoretical tracks at metallicities 0.2 $Z_{\odot}$, 1.0 $Z_{\odot}$, and 2.0 $Z_{\odot}$. In addition, the fraction of pixels assigned to shock ionization is highly dependent on the galaxy metal content. More metal-rich galaxies will tend to have their shock-ionized gas fraction underestimated, because their photoionization track is lower than that of a metal-poor galaxy. Because of the metallicity dependence of the line ratios in the diagnostic diagram, we maintain a metallicity grouping for the rest of our analysis. Specifically, we divide our sample into a sub-solar group, a solar group, and a super-solar group, as given by the galaxies’ metallicity. Within each constant metallicity group, the data components at different metallicity. For the [N II] diagnostics, we only have two sub-solar galaxies in terms of metal abundance, Holmberg II and NGC 4449. Tentatively, we put them into the sub-solar group. And, because the MSL is a theoretical limit for photoionization regardless of ionic species, we assume that the MSL for [N II] diagnostics is equivalent to the MSL of [S II] diagnostics for shock estimates.

4.2. [N II] versus [S II] Diagnostics

Figure 6 shows the theoretical tracks for the [N II] diagnostic (top) and the pixel diagnostic diagram for Holmberg II (middle) and NGC 4449 (bottom). The main difference between the [N II] diagnostic and the [S II] diagnostic is the amount of overlap between shock-ionization tracks and photoionization tracks. When we compare the two diagrams, we find that the [S II] diagnostic has smaller overlap between photoionization tracks and shock-ionization tracks than the [N II] diagnostic. The better shock discrimination offered by the [S II] diagnostic makes this diagram a preferred choice for this type of analysis. However, the shock tracks of the [S II] diagnostic at different metallicities overlap considerably, while the [N II] diagnostic shows a better separation among the tracks. To summarize, the [S II] diagnostic is effective at separating shocks from ionized gas, and the [N II] diagnostics is effective at separating gas components at different metallicity.

4.3. Single-line-ratio Diagnostics

When only a single line ratio is available due to limited observational conditions, this partial information still can be used for obtaining a rough measure of the metallicity or an approximate separation between shocks and photoionized components. The derived values will be less reliable than in the two-line-ratio diagnostics, due to degeneracies of the physical quantities in the single-line-ratio case. Figure 7 shows the theoretical tracks for each line ratio versus the normalized He II flux. For the [O III] ($\lambda 5007$)/H$\beta$ ratio, the shock tracks are well separated from the photoionization tracks for $Z = 2 Z_{\odot}$, but they tend to overlap for lower metallicity values ($Z \lesssim 1.0 Z_{\odot}$); rather, the ratio can be used as a metallicity indicator.

The [S II] ($\lambda \lambda 6716, 6731$)/H$\alpha$ ratio can be used to separate shocks from photoionized gas, e.g., by employing boundaries such as log([S II] ($\lambda \lambda 6716, 6731$)/H$\alpha$) > −0.5 for shocks. The [N II] ($\lambda 6583$)/H$\alpha$ ratio has a similar trend as the [S II] ($\lambda \lambda 6716, 6731$)/H$\alpha$ ratio, but shows a larger degree of degeneracy than the [S II] ($\lambda \lambda 6716, 6731$)/H$\alpha$ ratio, implying that it is a less effective single-line shock identifier.
Figure 5. [S ii] diagnostic diagram for each galaxy with theoretical grids. The gray lines show the theoretical grids presented in Figure 4. The error bars are presented at MUR spots: thin (cyan) error bars for 5σ detections and thick (blue) error bars for 10σ detections.

(A color version of this figure is available in the online journal.)

We should note that large [N ii]/Hα and [S ii]/Hα ratios can be obtained also in the presence of hot, diffuse photoionized gas (Reynolds et al. 1999, 2001). This gas will be characterized by weak [O iii]/Hβ ratios. In the absence of this line ratio, hot photoionized gas can be discriminated from shock-ionized gas by the morphological differences: the photoionized gas tends to be more uniformly distributed, while the shock-ionized gas is present in thin, shell-like regions. Overall, the use of a single line ratio as a shock diagnostic will produce an overestimate in the amount of shocks present in a region.

Figure 8 shows the observed pixel line ratio [O iii] (λ5007)/Hβ as a function of normalized Hα surface brightness for our galaxies. The metallicity dependence of the [O iii] (λ5007)/Hβ ratio at the bright end of the Hα flux follows...
Figure 6. Theoretical grids for the [N ii] diagnostics like Figure 4 (top). We plot all the models with the three metallicities, $Z = 0.2, 0.4, 1.0 Z_{\odot}$, in this single panel. The same error bars as presented in Figure 5. Diagnostic diagrams for Holmberg II (middle) and NGC 4449 (bottom) with theoretical tracks presented on the top panel. The [N ii] diagnostics shows the clear separation between different metallicities, but more overlaps between shock- and photoionization grids than the [S ii] diagnostics.

(A color version of this figure is available in the online journal.)

Figure 7. Line ratio of [O iii] $\lambda 5007$/H$\beta$, [S ii] $\lambda\lambda 6716, 6731$/H$\alpha$, [N ii] $\lambda 6583$/H$\alpha$ as a function of the normalized H$\alpha$ from K01 and A08. The models are the same ones presented in Figures 4 and 6. We plot the grids with different colors for each metallicity: $Z = 0.4 Z_{\odot}$ in green, $Z = 1.0 Z_{\odot}$ in blue, and $Z = 2.0 Z_{\odot}$ in magenta. The N350 photoionization models are located on the right-hand side of the N10 models due to their higher density and hence stronger H$\alpha$ emission. The shock-ionized gas has generally higher line ratios than photoionized gas. This property can be used as a shock separator, using $\log([S\ ii]^{\alpha}/H\alpha) > -0.5$. We can also find that the $[O\ iii]^{\lambda 5007}$/H$\beta$ ratios in H$\beta$ bright regions show a strong metallicity dependence as previously shown in Figure 4.

(A color version of this figure is available in the online journal.)

obtained for a larger number of pixels than in the case of two-line ratios, as the line images are deeper in the red than in the blue, but provide a poorer separating shock-/photoionization power than the diagnostic diagram. Internal variations of the...
Figure 8. Line ratio of $[\text{O} \text{ III}]$ ($\lambda$5007) / $\text{H} \beta$ as a function of the normalized $\text{H} \alpha$ surface brightness with the error bars presented in Figure 5. The surface brightness is normalized to the half-light radius surface brightness $\langle \Sigma \rangle$, i.e., the sum of all pixels over this value is equal to half of the total $\text{H} \alpha$ flux. The metallicity dependence of $[\text{O} \text{ III}]$ ($\lambda$5007) / $\text{H} \beta$ ratios in the bright end of $\text{H} \alpha$ flux fits well with the theoretical grids shown in the top panel of Figure 7. We can recognize the three groups from the figures: the sub-solar group (NGC 1569, NGC 5253, NGC 4449, Holmberg II, NGC 4214), the solar group (He 2–10, NGC 3077), and the super-solar groups (NGC 5236). Due to the dust lanes in NGC 5236, the horizontal pixel distribution near $[\text{O} \text{ III}]$ ($\lambda$5007) / $\text{H} \beta$ ≈ 0.0 is a propagated error from dust extinction correction. (A color version of this figure is available in the online journal.)

4.4. Relations between Shock-ionized Gas and Stellar Feedback

In this section, we will present our main results on the shock-ionized gas and its relations with the $H$-band magnitude (a proxy of stellar mass) and star formation (and star formation density). Using the MSL within each metallicity sub-sample, we separate pixels dominated by the shock ionization from those dominated by photoionization on the diagnostic diagram. Figure 10 shows the physical location of the identified shocks, re-projected onto the $\text{H} \alpha$ map of each galaxy. The pixels dominated by shock ionization are generally scattered in the outer rim of the central starburst.

From the identified shock-dominated pixels, we measure the two quantities. The first is the $\text{H} \alpha$ luminosity of the shocked component. If the shock is fully radiative, the mechanical luminosity driving the shock is about 20–80 times larger than the shocked $\text{H} \alpha$ luminosity (Rich et al. 2010). The numbers are dependent on the specific model, but generally we can adopt that a couple of percent of the total shock luminosity is radiated away through the $\text{H} \alpha$ emission. Hence, by measuring the $\text{H} \alpha$ luminosity of the shocks, we can estimate the underlying mechanical luminosity that is radiated away. The second is the...
ratio of the Hα luminosity of the shocked component to the total Hα luminosity. The total Hα luminosity is linked to the current SFR of the starburst. By adopting some appropriate history of star formation, we can calculate the mechanical luminosity from the SFR. Therefore, the ratio between the shocked and total Hα luminosity is an indicator of the balance between the energy injection rate and the energy loss rate.

When measuring the shocked Hα luminosity, we have three major factors affecting such a luminosity. The first one, already discussed in detail in previous sections, is the shock-/photoionization separation method, which is affected by a strong degeneracy linked to the presence of a continuum between the two gas components, rather than an abrupt transition. Within this framework, we suggest that the most important aspect for shock separation is consistency, rather than accuracy.

When we apply the MSL to our sample, the estimated shocks are consistent within each metallicity sub-sample due to the similar distribution on the diagnostic diagram, i.e., similar cooling pattern due to similar metallicity.

The second factor is a threshold driven by the observational detection limit. We adopt a 5σ threshold for each line emission. This is, however, an artificial threshold imposed by the depth of each image. In Section 4.4.2, we will apply a physical threshold for better consistency among different galaxies, although it will turn out that the qualitative results do not change if an artificial or physical threshold is applied.

The third factor is dust correction. As previously discussed, the dust correction on a pixel-by-pixel basis can suffer from bias due to stochastic uncertainty. But for heavily dust-obscured galaxies, the correction is necessary though the error can...
Figure 10. Distribution of identified shocks over the H$\alpha$ images (painted in green for sub-solar sample, blue for solar sample, and magenta for super-solar sample). Most identified shocks are located in the outer rim from the central or most prominent starburst. (A color version of this figure is available in the online journal.)

propagate to the shock estimates. In our sub-solar sample, only NGC 1569 and NGC 4449 have some changes from this correction. Like the detection threshold issue, the dust correction does not change the qualitative results.

Finally, two additional sources of confusion can lead to either the underestimate or the overestimate of the shock luminosity. The first is the presence of shocks in strongly photoionized regions; in this case, the shock emission will be diluted, and we will tend to underestimate the shock luminosity. The second is the presence of hot, diffuse photoionized gas; in this case, the low-ionization lines will be enhanced, and, especially if utilizing a single-line-ratio diagnostic (e.g., [S\textsc{ii}]/H$\alpha$), the diffuse photoionized gas may be mistaken for shock emission and the shock luminosity overestimated. The latter scenario can, however, be controlled by investigating the morphology of the high [S\textsc{ii}]/H$\alpha$ regions: photoionized gas will tend to be diffuse, while shocks will tend to present a filamentary morphology. In what follows, we assume that these two additional sources of bias are small and roughly compensate each other, when galaxy-integrated properties are investigated.

4.4.1. Correlations between $L_{\text{shock}}$, $L_{\text{shock}}/L_{\text{tot}}$, $\Sigma_{\text{SFR,HL}}$, and $M_H$

Tables 3 and 4 report various quantities derived from the extinction-corrected H$\alpha$ flux and luminosity of each galaxy. These fluxes and luminosities are derived from the line emission images, after imposing a 5$\sigma$ threshold to each line image, including those lines used for the diagnostic diagrams. While line fluxes and luminosities are corrected for the effects of dust extinction, using the H$\alpha$/H$\beta$ line ratio, the diagnostic line ratios are used without extinction corrections. When $E(B-V)$ is smaller than 0, we assign $E(B-V) = 0$.

We measure two kinds of H$\alpha$ fluxes and luminosities: $F_{\text{H}\alpha,\text{tot}}$ (Table 3, Column 2), which is the H$\alpha$ flux derived from the whole H$\alpha$ image, and $F_{\text{H}\alpha,\text{tot,diag}}$ (Table 3, Column 4), which is the H$\alpha$ flux summed over all pixels above the 5$\sigma$ detection limit. This second definition mirrors the diagnostic diagram.
Figure 11 summarizes the results listed in Table 4, with the sub-solar sample in green, the solar sample in blue, and the super-solar sample in magenta. Though many parameters seem to have no significant relation with one another, we find three suggestive correlations in the sub-solar group, marked in the figure with green dotted lines. The green lines are $\chi^2$ minimization fits to the most refined data presented in the following subsection (Figure 13) after additional corrections:

$$\log(L_{\text{shock}}) = 0.62(\pm 0.05) \times \log(\text{SFR}) + 39.9(\pm 0.05)$$

$$= 0.92(\pm 0.41) \times \log(\Sigma_{\text{SFR, HL}}) + 38.8(\pm 0.26),$$

$$\log(L_{\text{shock}} / L_{\text{tot}}) = 0.26(\pm 0.08) \times M_H + 4.25(\pm 1.58)$$

$$= -0.65(\pm 0.2) \times \log(L_H / L_{H, \odot}) + 5.11(\pm 1.60),$$

where $L_{\text{tot}}$ is the total H$\alpha$ luminosity (second column of Table 3), $L_{\text{shock}}$ is the total H$\alpha$ luminosity from the shock-ionized gas component (second column of Table 4), $\Sigma_{\text{SFR, HL}}$ is the SFR density using the half-light pixel area (fifth column of Table 4), and $L_H$ is an H-band luminosity converted from $M_H$ by adopting $M_{H, \odot} = 3.32$ (see the next section for the details about additional corrections and correlation tests for the equations above). Here we focus on the qualitative interpretation, which remains unchanged even after the application of refined corrections. For convenience, we drop the “H$\alpha$” subscripts for the H$\alpha$ luminosities used in Tables 3 and 4. It is interesting that all the “diag” quantities show worse correlations with $L_{\text{shock}}$ than the other quantities. This implies that the artificial detection threshold, which affects both the $L_{\text{shock}}$ and “diag” quantities, smooths out the underlying physical relations. Therefore, we exclude all the “diag” quantities hereafter, and $L_{\text{shock}}$ is the only quantity derived from the diagnostic diagram among the four quantities, $L_{\text{shock}}$, $\Sigma_{\text{SFR, HL}}$, and $M_H$ in Equations (3)–(6). To reiterate, only $L_{\text{shock}}$ is a biased measurement driven

| Galaxy  | $L_{\text{H$\alpha$, sh}}$ (erg s$^{-1}$ cm$^{-2}$) | SFR$^b$ (M$_{\odot}$ yr$^{-1}$) | SFR$_{\text{diag}}$ (M$_{\odot}$ yr$^{-1}$) | $\Sigma_{\text{SFR, HL}}$ (M$_{\odot}$ yr$^{-1}$ kpc$^{-2}$) | $\Sigma_{\text{diag}}$ (M$_{\odot}$ yr$^{-1}$ kpc$^{-2}$) | $L_{\text{H$\alpha$, sh}} / L_{\text{H$\alpha$, tot}}$ | $L_{\text{H$\alpha$, sh}} / L_{\text{H$\alpha$, diag}}$ | $A_{\text{sh}} / A_{\text{tot, diag}}$ |
|---------|------------------------------------|-------------------------------|-------------------------------|------------------------------------|------------------------------------|----------------------------------|-------------------------------|-------------------------------|
| NGC 1569 | $3.29 \times 10^{39}$ | 0.13 | 0.067 | 9.6 | 1.5 | 0.20 | 0.26 | 0.28 |
| NGC 5253 | $4.69 \times 10^{39}$ | 0.26 | 0.11 | 10 | 0.35 | 0.14 | 0.15 | 0.38 |
| NGC 4214 | $1.86 \times 10^{39}$ | 0.09 | 0.067 | 2.3 | 0.26 | 0.16 | 0.18 | 0.46 |
| He 2–10 | $1.74 \times 10^{39}$ | 0.62 | 0.56 | 33 | 1.4 | 0.022 | 0.024 | 0.28 |
| NGC 3077 | $2.93 \times 10^{38}$ | 0.073 | 0.052 | 2.6 | 0.034 | 0.048 | 0.21 |
| NGC 5236 | $1.90 \times 10^{38}$ | 1.52 | 0.18 | 100 > | 2.3 | 0.097 | 0.43 | 0.54 |
| Holmberg 2 | $7.20 \times 10^{38}$ | 0.019 | 0.0032 | 1.4 | 0.20 | 0.29 | 0.64 | 0.54 |
| NGC 4449 | $6.86 \times 10^{38}$ | 0.58 | 0.25 | 4.6 | 0.67 | 0.094 | 0.18 | 0.18 |

Notes.

$^a$ Total H$\alpha$ luminosity with dust extinction correction, associated with shock-ionized gas.

$^b$ SFR derived from the total H$\alpha$ luminosity in the second column of Table 3.

$^c$ SFR derived from the H$\alpha$ luminosity in the diagnostic diagram. We subtract the shock-ionized component too for the H$\alpha$ luminosity.

$^d$ SFR density defined as $\Sigma_{\text{SFR, HL}} = SFR / A_{\text{SFR, HL}}$, where the half-light area, $A_{\text{SFR, HL}}$, is the pixel area above the half-light surface brightness. For NGC 5236, we use the A1 section from H11. Because the section covers the bright central region, the half-light flux is higher than the level from the entire image. Hence, the $A_{\text{SFR, HL}}$ for NGC 5236 is underestimated. The estimated SFR density is several hundreds.

$^e$ SFR density defined as $\Sigma_{\text{diag}} = SFR_{\text{diag}} / A_{\text{tot, diag}}$, where the $A_{\text{tot, diag}}$ is the pixel area of all data points in the diagnostic diagram.

$^f$ Luminosity ratio to the total H$\alpha$ luminosity.

$^g$ Luminosity ratio to the total H$\alpha$ luminosity in the diagnostic diagram.

$^h$ Area ratio of the shock-ionized component to the area of all the pixels on the diagnostics.
by the detection threshold and by the adopted method for separating shocks from photoionized gas in the diagnostic diagram.

Since we emphasize consistency over absolute accuracy, we derive our two main, tantalizing qualitative results:

1. \(L_{\text{shock}} \) increases with SFR and \(\Sigma_{\text{SFR,HL}}\); and
2. \(L_{\text{shock}}/L_{\text{tot}}\) increases as \(-M_{H}\) decreases.

The first result implies that a larger energy injection from a higher SFR drives stronger radiative shocks into the surrounding ISM. This is a quite intuitive result that we can generally expect. The interesting point is the sub-linear (i.e., the slope in Equation (3) is lower than unity) relation between \(L_{\text{shock}}\) and SFR, implying that the efficiency driving radiative shocks seems not to increase as much as SFR increases. Another interesting point is that we find the relation between \(\Sigma_{\text{SFR,HL}}\) and \(L_{\text{shock}}\) too. The half-light area is a quantity representing the compactness of a star-forming region. \(\Sigma_{\text{SFR,HL}}\) will increase for the same SFR if the star-forming region is more compact. This implies that, even for an identical amount of total energy injection, \(L_{\text{shock}}\) can be larger if the injected energy is deposited in a smaller volume. In addition, we observe a stronger distinction between different metallicity subgroups in the \(L_{\text{shock}}\) versus \(\Sigma_{\text{SFR,HL}}\) relation than in the \(L_{\text{shock}}\) versus SFR relation (top left panel of Figure 13). This is due to (1) a systematic underestimate of \(L_{\text{shock}}\) due to the lower \([\text{O iii}]\) \((\lambda 5007)/H\beta\) ratios for more metal-abundant galaxies and (2) more compact star-forming morphology for more massive galaxies (hence, higher metallicity due to mass–metallicity relation). This distinction is shown by the green, blue, and magenta dashed lines in Figure 13.

The second result shows that \(L_{\text{shock}}/L_{\text{tot}}\) becomes smaller for brighter galaxies in the \(H\) band. If we consider that the gravitational potential well at the center of galaxies is mostly shaped by the stellar mass, for which \(M_{H}\) is a good approximation, the result shows the effect of the gravitational potential well on the strength of feedback-driven shocks. Higher \(L_{\text{shock}}\) will be present in a shallower potential well, because the injected energy in a shallower potential well can be transported out more easily than in a deeper potential well.

Though the quantitative values will be improved in the next section, the qualitative interpretations can be summarized...
Figure 12. Relations between SFR–$L_{\text{sh}}$ (top left), SFR–$L_{\text{sh}}/L_{\text{tot}}$ (top right), $M_{\text{HI}}$–SFR (bottom right), and $M_{\text{HI}}$–$L_{\text{sh}}/L_{\text{tot}}$(bottom right), presented in Figures 1 and 11. Due to SFR $= \kappa L_{\text{gas}}$, the two top panels are equivalent. The lines represent the theoretical ratio $\mu/(1 + \mu)$, where $\mu = \kappa L_{\text{sh}}/\lambda$ in Equations (19) and (20). The range of $\mu$ can cover the observed data. Because $M_{\text{HI}}$–SFR is not independent (Noeske et al. 2007), the relations in the right panels are projected results from the underlying three-dimensional relation of $L_{\text{sh}}/L_{\text{tot}}$– $M_{\text{HI}}$. The bottom right panel of Figure 13 shows the preliminary three-dimensional relation.

(A color version of this figure is available in the online journal.)

\[\text{As } (1) \quad L_{\text{shock}} \propto \text{SFR}^{0.62} (\text{or } L_{\text{shock}}/L_{\text{tot}} \propto \text{SFR}^{–0.38}), \quad (2) \quad L_{\text{shock}} \propto \text{SFR}^{0.92}, \quad \text{and } (3) \quad L_{\text{shock}}/L_{\text{tot}} \propto 10^{0.26} \log H_{\text{II}} \quad (\text{or } L_{\text{shock}}/L_{\text{tot}} \propto (L_{\text{H}}/L_{\text{HI,} \odot})^{–0.65}).\]

4.4.2. The Impact of Detection Thresholds and Dust Extinction

To investigate our results in a quantitative way, we discuss the impact of detection thresholds and dust extinction corrections on the estimate of $L_{\text{shock}}$. It will turn out that this refining process for investigation yields better support for our results.

Detection thresholds. The results shown in the previous section are based on $5\sigma$ detection limits for the emission line images. As a test, we apply two additional thresholds, $3\sigma$ and $7\sigma$, to verify how much changing the detection threshold changes $L_{\text{shock}}$. In general, lowering the detection threshold will admit more pixels in the diagnostic diagram and increase the H$\alpha$ luminosity of the shock; the opposite happens for higher detection thresholds.

The effect of changing the threshold is shown in Figure 13, where the $L_{\text{shock}}$ estimates for the $3\sigma$, $5\sigma$, and $7\sigma$ cuts are reported as red points connected by vertical bars. This demonstrates that the depth of the detection cut significantly affects the estimates of $L_{\text{shock}}$. $L_{\text{shock}}$ is two times larger for the $3\sigma$ threshold than for the $7\sigma$ threshold. This is because the shock-ionized gas is generally fainter and more diffuse than the photoionized gas (see CO4 and H11), and the faint pixels are those that are more prominently recovered by a lower threshold. Our result suggests the existence of non-negligible amounts of shocked gas below our detection limit and emphasizes the need for an internally consistent measurement over an accurate measurement.

We now apply a physical threshold to the emission line images, in order to check the effect of the bias caused by the different exposure depths of each galaxy. We design the physical threshold by imposing a cut on the absolute surface brightness:

\[I(\lambda)_{\text{cut}} = 10^{-15} \text{erg cm}^{-2} \text{arcsec}^{-2}, \quad (7)\]

\[F(\lambda)_{\text{cut}} = I(\lambda)_{\text{cut}} 10^{-0.4(E(B-V)_{\text{shock}} + k_{\lambda} \Omega_{\text{bin}})}, \quad (8)\]

with Equation (7) converting the surface brightness to a total flux, $E(B-V)_{\text{gas}}$ is the foreground color excess of the gas, $k_{\lambda}$ is the Milky Way extinction curve, and $\Omega_{\text{bin}}$ is the solid angle subtended by our pixel bin size. We set the above surface brightness threshold to the H$\alpha$, H$\beta$, [O ii] (15007), and [S ii] (\lambda 6716, 6731) images. For [N ii] (6583), we use the relations [N ii] (\lambda 6548, 6583) $\approx 0.5 \times$ [S ii] (\lambda 6716, 6731) and [N ii] (\lambda 6548, 6583) $\approx 1.3 \times$ [N ii] (6583) for its equivalent brightness cut. We apply this threshold to all the galaxies in our sub-solar sample. This even physical threshold also can reduce the systematic bias from the MUR spot effect presented before.

The green points in Figure 13 show the effect of imposing the absolute surface brightness cut. For NGC 4214, NGC 4449, NGC 5253, and Holmberg II, the physical threshold is equivalent to the $5\sigma$–$7\sigma$ detection limit. For NGC 1569, on the other hand, the images are relatively shallow due to its unusually high foreground extinction, $E(B-V) = 0.70$. The physical threshold for NGC 1569 corresponds to a $1\sigma$–$2\sigma$ detection limit, which we do not apply to the data, as far too shallow a limit. However, we report this result for NGC 1569 as an upward-pointing green arrow in Figure 13, to remark that our estimates are lower limits to the actual $L_{\text{shock}}$.

Dust extinction. In general, the extinction vectors on the diagnostic diagram move the line ratios into the photoionized area. Therefore, the dust-corrected estimates of $L_{\text{shock}}$ decrease after dust extinction correction. Since the extinction vector is...
Figure 13. Corrections to $L_{\text{shock}}$: apparent thresholds, 3σ, 5σ, and 7σ (red bar), absolute threshold (green point), and dust correction to the line ratios (blue point). The dashed lines represent $L_{\text{shock}} \propto \text{SFR}^{0.62 \pm 0.05}$, $L_{\text{shock}} \propto \Sigma_{\text{SFR,II}}^{0.92 \pm 0.41}$, and $L_{\text{shock}}/L_{\Sigma} \propto 10^{0.26 \pm 0.08} M_\odot$ (or $\propto (L_{\text{H}}/L_{\text{tot}})^{0.05 \pm 0.2}$). Holmberg II and NGC 4449 are marked with an asterisk (*) to indicate their uses of the [NII] diagnostics, while the others of the [SII] diagnostics. From the red bars, we can find that a non-negligible number of shocks can be found from a deeper observation. NGC 1569 is indicated as an arrow because the absolute cut corresponds to $1\sigma$--$2\sigma$ levels for the line emission images. Due to the heavy dust extinctions of NGC 1569, the dust correction should be significant for NGC 1569. We mark the dust correction effect as a blue arrow, which works in the opposite direction of the absolute cut effect. NGC 4449 and NGC 5253 turn out to have significant dust corrections (see the green and blue points). Though we have limited number statistics, this figure shows strong evidence for the existence of regulations between the four quantities $L_{\text{shock}}$, $L_{\text{shock}}/L_{\Sigma}$, SFR, and $M_\star$. The bottom right panel shows the three-dimensional plot of $L_{\text{shock}}/L_{\Sigma}$ vs. $M_\star$ and SFR. The four blue points are the same blue points on the other panels for NGC 4449, NGC 5253, NGC 4214, and Holmberg II. NGC 1569 is plotted in red with its 5σ cut result. The sub-linear scaling of $L_{\text{shock}} \propto \text{SFR}^{0.62 \pm 0.05}$ can be due to projection effects of the three-dimensional relation.
(A color version of this figure is available in the online journal.)

small for the diagnostic diagram, the dust extinction corrections generally produce minor effects on the line ratios. But if many data points are crowded near our shock separating line, even the small extinction vector can migrate a non-negligible number of pixels from the shock-ionized into the photo ionized area. We find that NGC 4449 and NGC 1569 have relatively high dust corrections and NGC 5253 has a non-negligible dust correction.

Figure 14 shows the distribution of $E(B-V)$ color excesses used for each galaxy on a pixel-by-pixel basis. The right panel shows the absolute counts of the shock-ionized bins for dust correction, and the left panel shows the normalized counts. The foreground extinction values are given with the names in parentheses. As mentioned previously, NGC 1569 suffers from heavy foreground dust extinction, and hence it shows the most significant color excess. On the other hand, NGC 4449 and NGC 5253 have relatively high internal dust extinctions, while their foreground dust extinctions are much smaller than NGC 1569: $E(B-V) = 0.019$ for NGC 4449 and 0.056 for NGC 5253. The other two galaxies, Holmberg II and NGC 4214, show minor dust extinctions. We can expect that the dust correction will affect the shock estimates for NGC 4449, NGC 5253, and NGC 1569.

The blue points in Figure 13 show the estimates of shocks after dust corrections are applied to the diagnostic diagrams. We can readily observe that the dust corrections on the diagnostic diagrams have some effect on NGC 4449 and NGC 5253. For NGC 1569, we mark the dust correction effect as a blue arrow, since the dust correction will work in the opposite direction (reducing the shock estimate) of the physical threshold. And the correction should be the most for NGC 1569. For NGC 4214 and Holmberg II, the dust extinctions are small. Both galaxies are relatively free from all of the issues related to dust extinction correction. NGC 4449 and NGC 5253 show a similar color excess in the left panel (for the normalized count of bins) of Figure 14. This shows that they have a similar dust extinction property. However, NGC 4449 is brighter than NGC 5253 in Hα emission; hence, the dust correction affects more the estimates of bins containing shocked gas in NGC 4449 than in NGC 5253.
Section 4.4.2, the dust correction affects proportionally NGC 4449 more than NGC 5253 to its heavy foreground extinction. The dust extinction corrections are not negligible in NGC 5253, NGC 4449, and NGC 1569. NGC 5253 and NGC 4449 show significant correlation.

Because no exact correction is available to NGC 1569 due to shallower observation, we take the values between 0 and 1. A smaller value of significance indicates a more robust against these effects.

Notes.

Table 5

| Relation     | Data | Pearson | Spearman | Kendall |
|--------------|------|---------|----------|---------|
| $L_{\text{sh}}$ vs. SFR | −NGC 1569 | 0.994 | 1.000 (0.000) | 1.000 (0.042) |
| +NGC 1569 | 0.963 | 1.000 (0.000) | 1.000 (0.014) |
| $L_{\text{sh}}$ vs. $\Sigma_{\text{HI}}$ | −NGC 1569 | 0.844 | 0.800 (0.200) | 0.667 (0.174) |
| +NGC 1569 | 0.829 | 0.700 (0.188) | 0.600 (0.142) |
| $L_{\text{sh}}/L_{\text{tot}}$ vs. $M_{\text{HI}}$ | −NGC 1569 | 0.917 | 1.000 (0.000) | 1.000 (0.042) |
| +NGC 1569 | 0.898 | 0.900 (0.037) | 0.800 (0.050) |

The Astrophysical Journal, 777:63 (21pp), 2013 November 1

HONG ET AL.

5. DISCUSSION

5.1. Interpretation of $L_{\text{shock}}/L_{\text{tot}}$

The deposited energy from stellar feedback, $E_{\text{mech}}$, cools away by various cooling mechanisms, $E_{\text{loss}}$. The remaining energy, $E_{\text{wind}}$, can drive galactic scale winds

$$E_{\text{wind}}(t) = E_{\text{mech}}(t) - E_{\text{loss}}(t).$$

The amount of feedback energy deposited into the surrounding ISM (and its rate) can be calculated from stellar population synthesis models (STARBURST99 in this paper) when we have (or assume) a star formation history, $h_*(t)$,

$$E_{\text{mech}}(t) = \int_{-\infty}^{t} L_{\text{mech}}(t') dt',$$

$$L_{\text{mech}}(t) = \int_{-\infty}^{t} K_{\text{inst}}(t-t') h_*(t') dt',$$

where $K_{\text{inst}}(t)$ is a kernel function of luminosity evolution for instantaneous star formation. Many models and assumptions are involved in calculating the luminosity, such as the metallicity, initial mass function (IMF), stellar atmosphere models, and stellar evolution tracks. We need to keep in mind this complexity when modeling and interpreting stellar feedback.

The two models most commonly used are the instantaneous burst, $h_*(t) = \delta(t - 0)$, and the continuous star formation, $h_* = \pi(t - 0)$, here written mathematically using the conventional delta function and step function. The mechanical luminosity from stellar feedback for each star-forming model can be written as

$$L_{\text{inst}}(t) = \delta(t - 0),$$

$$L_{\text{cont}}(t) = \pi(t - 0),$$

$$\approx \pi, K_0(t > 40 \text{Myr}).$$

The two luminosity templates, $K_{\text{inst}}(t)$ and $K_{\text{cont}}(t)$, can be provided by STARBURST99 (Leitherer et al. 1999). One of the important results from the templates is that, after 40 Myr, the luminosity of the continuous star formation model is stabilized to a constant, $K_0 \approx 10^{32} \text{erg s}^{-1}$, while the luminosity of the instantaneous model fades out to $<10^{26} \text{erg s}^{-1}$, because all massive stars explode as SNe within 40 Myr.
Because the hydrogen recombination lines are tracers of SFR in a very short term period (<10 Myr; Leitherer et al. 1999), the obtained SFR can be considered as an instantaneous measure of current star formation. Ideally, we have to subtract $L_{\text{shock}}$ from $L_{\text{tot}}$ to obtain $h_*$ (as a reminder, $L_{\text{shock}}$ and $L_{\text{tot}}$ are Hα luminosities). Because $L_{\text{shock}}$ is a small fraction of $L_{\text{tot}}$ for most normal star-forming galaxies and is generally hard to measure $L_{\text{shock}}$, as explained throughout this paper, conventionally $L_{\text{shock}}$ has been ignored:

$$h_* = \kappa (L_{\text{tot}} - L_{\text{shock}})$$

$$\approx \kappa L_{\text{tot}}; \quad \kappa = 7.9 \times 10^{-42}.$$  \hspace{1cm} (15)

In our sample, the fraction of $L_{\text{shock}}$ is at most 0.3, so the conventional approximation would be acceptable also in our case. We keep, however, the explicit formula, Equation (15), which will be used later to derive $L_{\text{shock}}/L_{\text{tot}}$. From the observed SFR, $h_*$, we can rewrite the mechanical luminosity of the continuous burst model after 40 Myr, $L_{\text{mech}} = \kappa (L_{\text{tot}} - L_{\text{shock}}) K_0 \approx \kappa L_{\text{tot}} K_0 (t > 40 \text{ Myr})$.  \hspace{1cm} (16)

In general, $K_0$ can be used in place of the time-dependent templates, $K_{\text{tot}}(t)$ or $K_{\text{con}}(t)$.

Gas cooling is more complicated to estimate than stellar feedback energy because thermodynamic and hydrodynamic interactions are involved between surrounding gas and feedback energy. In the early stages of star formation, the energy deposited into the ISM produces a hot bubble that expands adiabatically (Weaver et al. 1977). Hence, most of the feedback energy is stored as $E_{\text{wind}}$ and is invisible in the optical. As the hot bubble adiabatically cools, the expanding shock front becomes radiative. Our measured $L_{\text{shock}}$ is a tracer of this radiative shock. The total radiative loss, $L_{\text{shock loss}}$, can be estimated from the observed Hα loss, $L_{\text{shock}}$,

$$L_{\text{shock loss}} = \lambda L_{\text{shock}}.$$  \hspace{1cm} (18)

The conversion factor, $\lambda$, depends on the shock models and ISM properties, such as metallicity, ambient gas density, magnetic field strength, and pre-ionization fraction. Rich et al. (2010) estimate $\lambda = 20–80$ for slow radiative shocks. We use this range of values for our qualitative interpretation, but note the complexity of treating radiative shocks, which includes the uncertainties in the parameters chosen for the population synthesis models.

Now we derive an estimate of $L_{\text{shock}}/L_{\text{tot}}$ from Equations (17) and (18):

$$\frac{L_{\text{shock}}}{L_{\text{tot}}} = \frac{\mu}{1 + \mu} \approx \mu,$$  \hspace{1cm} (19)

$$\mu \equiv \left(\frac{K_0}{\lambda}\right) \left(\frac{L_{\text{shock}}}{L_{\text{mech}}}\right).$$  \hspace{1cm} (20)

The term $\mu/(1 + \mu)$ is the explicit derivation from $h_* = \kappa (L_{\text{tot}} - L_{\text{shock}})$ and can be approximated to $\mu$ when $\mu$ is small enough, $\mu < 0.1$. Our observed ratios, $L_{\text{shock}}/L_{\text{tot}} = 0.05–0.40$, justify the approximation, as we discuss below. When we take the fiducial values, $K_0 = 10^{42}$ and $\lambda = 60$, we obtain ($K_0/\lambda) = 0.13$. If we allow a larger range of values to cover most physical conditions, $K(t) = 10^{36–42}$ and $\lambda = 20–80$, we obtain ($K_0/\lambda) = 0.00–0.40$, which is in agreement with the observed range of ratios for $L_{\text{shock}}/L_{\text{tot}} = 0.05–0.40$. Indeed, as shown in Figure 12, the observed range and theoretical expectations overlap, implying that in our sample $L_{\text{shock loss}}/L_{\text{mech}} \approx 1$.

In order to attempt an explanation of the scaling relation $L_{\text{shock}}/L_{\text{tot}} \propto (L_H/L_{\odot,H})^{0.65}$, we discuss the three parameters $K_0, \lambda$, and $L_{\text{shock loss}}/L_{\text{mech}}$ in greater detail. $\kappa$ is dictated by atomic physics and will not be discussed further. As presented above, $K_0$ depends on the IMF, stellar atmosphere models, stellar evolution models, stellar metallicity, and star formation history, while $\lambda$ is a function of the shock model and the ISM properties. As we have divided our sample into sub-samples according to metal content, we can assume that variations in stellar metallicity and ISM properties are minimized. Furthermore, IMF, stellar atmosphere models, and evolution models are not (likely) a function of the galactic environment. The star formation history of our galaxies has remained constant over the past few tens of Myr at least; thus, $K_0$, is likely to have remained roughly constant. If we assume that the dependency of $\lambda$ over the specific shock model is small, then the role of ($\kappa K_0/\lambda$) in Equation (20) is to simply set the absolute scale of the relation. Much of the environment dependency, i.e., the dependency on $H$, is carried by $L_{\text{shock loss}}/L_{\text{mech}}$.

The ratio $L_{\text{shock loss}}/L_{\text{mech}}$ represents the energy balance between the gain $L_{\text{mech}}$ from feedback energy and loss $L_{\text{shock loss}}$ by radiative shock. The higher the value, the higher the loss of feedback energy through radiation. To drive a gas outflow that may develop into a galactic superwind, the hot bubble should retain sufficient kinetic energy to expand to a scale height comparable to or larger than that of the galactic disk (e.g., MacLow et al. 1989; de Young & Heckman 1994; Murray et al. 2011). For our galaxies, we have a high likelihood that no such major gas outflow will be driven out of the galaxy: we derive $L_{\text{shock loss}}/L_{\text{mech}} \approx 1$, although it only represents the current strength of underlying gas expansion driving radiative shocks. The case of $L_{\text{shock loss}}/L_{\text{mech}} > 1$ is very unlikely, when we include other channels of energy loss, such as X-ray emissions. But still it is not impossible since a lot of energy can be cumulated during the early adiabatic phase and then released through radiative shocks in a short period of time. Overall, there is a strong argument for the environment (stellar mass) dependency in our scaling relation to be carried by $L_{\text{shock loss}}/L_{\text{mech}}$.

5.2. $L_{\text{shock}}/L_{\text{tot}} \propto (L_H/L_{\odot,H})^{-\alpha(SFR)}^{-\beta}$

Given our limited sample size and the radiative nature of the stellar feedback we observe, the content of this section is speculative at best, but provides some tantalizing suggestions and a direction for improvement and progress.

Our two scaling relations, (1) $L_{\text{shock}} \propto SFR^{0.62}$ and (2) $L_{\text{shock}}/L_{\text{tot}} \propto (L_H/L_{\odot,H})^{0.65}$, are suggestively similar to those derived in a recent simulation by Hopkins et al. (2012). These authors derive relations between the wind mass-loss rate driven by stellar feedback, $M_{\text{wind}}$, and the two quantities of SFR, $M_*$, and stellar mass, $M_*$: (1) $M_{\text{wind}} \propto M_*^{0.7}$ and (2) $M_{\text{wind}}/M_* \propto (M_*/M_\odot)^{-(0.25–0.5)}$. The scaling relations derived by Hopkins et al. are for feedback-driven galactic winds, while our relations are for radiative shocks, so the similarities should be taken with caution at this stage. However, if the way in which stellar feedback scales with both SFR and stellar mass is independent of the fate of the feedback energy, we can postulate that our observed scaling relations, derived for radiative shocks, reproduce those derived from simulations of end-phase galactic superwinds.
Conversely, also the opposite argument could be made. Our observational scaling relations indicate that radiative losses, \( L_{\text{shock}} / L_{\text{tot}} \), decrease for increasing stellar mass. In more massive galaxies the level of radiative losses is small so that most of the stellar feedback energy then is available to drive gas motions and, possibly, a wind. On the other hand, we find that radiative losses in dwarf galaxies are substantial and, thus, may reduce the efficiency of driving mass loss. The effects of radiative losses, therefore, go in an opposite sense to the depth of radiative losses in dwarf galaxies are substantial and, thus, may be identified as shock-dominated systematically decreases for increasing metallicity, indicating that shocks are underestimated in higher metallicity galaxies when the MSL is applied. A common problem to our analysis is that weak shocks projected on strongly photoionized regions will not be identified, likely leading to some underestimate of the total shock luminosity in a galaxy.

2. The [S ii] diagnostic is preferred for shock discrimination to the [N ii] diagnostic, because it has less overlap between the photoionized and the shock-ionized components. Conversely, the [N ii] diagnostics is more sensitive to metallicity than the [S ii] diagnostics.

3. The distribution of single-line ratios is consistent with theoretical expectations. Due to its metallicity dependence, the [O iii]/H\( \beta \) ratio provides poor discrimination between shocks and photoionized gas, especially for metal-poor galaxies. Both [N ii]/H\( \alpha \) and [S ii]/H\( \alpha \) can be used as single-line ratios for estimating shocks, with the [S ii]/H\( \alpha \) ratio providing a better shock discriminator, because of its lower degree of degeneracy for varying physical conditions. In general, two-line-ratio diagnostics should be preferred over single-line-ratio diagnostics, because the latter are sensitive not only to shocks but also to hot, photoionized gas, thus leading to an overestimate of the shock luminosity. However, morphology can help discriminate between the two cases, since photoionized gas is more uniformly diffused, while shock-ionized gas has a more filamentary structure.

4. We find that: (1) a larger H\( \alpha \) luminosity from shocks, \( L_{\text{shock}} \), is found for higher SFR with a sub-linear scaling; and (2) a higher ratio of the shock to the total H\( \alpha \) luminosity, \( L_{\text{shock}} / L_{\text{tot}} \), is obtained for increasing absolute H-band magnitudes, \( M_H \) (lower stellar masses). The two scaling relations are expressed as \( L_{\text{shock}} \propto SFR^{0.62} \) and \( L_{\text{shock}} / L_{\text{tot}} \propto (L_H / L_{\odot,H})^{-0.65} \). These results have been obtained for the sub-solar sample, NGC 1569, NGC 5253, NGC 4449, Holmberg II, and NGC 4214. No similar conclusions can be derived at this point for the other metallicity groups, due to the small sample statistics, although their trends are consistent with those of the sub-solar sample. We find tantalizing similarities between our observationally derived scaling relations and those recently obtained from simulations of galactic superwinds.

5. Our results are derived for starburst galaxies, i.e., for galaxies where we expect the effects of stellar feedback to have the most impact on the surrounding environment. Quiescently star-forming galaxies can be expected to have less substantial SN-powered winds.

Due to the limited sample statistics, our results need to be refined by larger samples, especially in the high-metallicity bins. However, we find evidence for the existence, in starburst galaxies, of regulation of the strength of radiative shocks by two galactic parameters: the SFR and the stellar mass.

6. SUMMARY

With HST narrow- and broad-band imaging, we have identified stellar-feedback-induced shocks in a sample of eight local starburst galaxies and investigated the relations among shock-related quantities, stellar mass, and star formation. We perform our analysis in a spatially resolved fashion, sampling our galaxies in pixels about 3–15 pc in size. Though our number statistics are still limited, we have found some indication of a correlation between the observed shock-ionized gas and the star-forming activity. We summarize below.

1. After dividing our sample into three subgroups according to their metallicities, sub-solar (NGC 1569, NGC 5253, NGC 4449, Holmberg II, NGC 4214), solar (He 2–10, NGC 3077), and super-solar (NGC 5236), we apply the “MSL” in an internally consistent manner to separate shock-ionized gas from photoionized gas in a classical diagnostic plot of the [O iii]/H\( \beta \) ratio versus the [S ii]/H\( \alpha \) (or [N ii]/H\( \alpha \)) ratio. The fraction of pixels identified as shock-dominated systematically decreases for...
contract NAS 5-26555. These observations are associated with programs 10522 and 11146. The paper also uses Early Release Science observations made by the WFC3 Science Oversight Committee (program 11360). We are grateful to the Director of STScI for awarding Director’s Discretionary time for this program.

REFERENCES

Allen, M. G., Groves, B. A., Dopita, M. A., Sutherland, R. S., & Kewley, L. J. 2008, ApJS, 178, 20 (A08)
Baldwin, J. A., Philips, M. M., & Terlevich, R. 1981, PASP, 93, 5
Berg, D. A., Skillman, E. D., Marble, A. R., et al. 2012, ApJ, 754, 98
Binette, L., Dopita, M. A., & Tuohy, I. R. 1985, ApJ, 297, 476
Bresolin, F. 2011, ApJ, 729, 56
Bresolin, F., Schaerer, D., González Delgado, R. M., & Stasinska, G. 2005, A&A, 441, 981
Calzetti, D. 2001, PASP, 113, 1449
Calzetti, D., Harris, J., Gallagher, J. S., et al. 2004, AJ, 127, 1405 (C04)
Cardelli, J. A., Clayton, G. C., & Mathis, J. S. 1989, ApJ, 345, 245
Cecil, G., Bland-Hawthorn, J., & Veilleux, S. 2002, ApJ, 576, 745
Chevalier, R. A., & Clegg, A. W. 1985, Natur, 317, 44
Croxall, K. V., van Zee, L., Lee, H., et al. 2009, ApJ, 705, 723
Davé, R., Oppenheimer, B. D., & Finlator, K. 2011, MNRAS, 415, 11
De Young, D. S., & Heckman, T. M. 1994, ApJ, 431, 598
Engelbracht, C. W., Rieke, G. H., Gordon, K. D., et al. 2008, ApJ, 678, 804
Fruchter, A., Sosey, M., Hack, W., et al. 2009, The MultiDrizzle Handbook, Version 3.0 (Baltimore, MD: STScI)
Gonzaga, S., & Biretta, J. 2010, in HST WFPC2 Data Handbook, Version 5.0 (Baltimore, MD: STScI)
Guseva, N. G., Izotov, Y. I., Stasińska, G., et al. 2011, A&A, 529, A149
Heckman, T. M., Armus, L., & Miley, G. K. 1990, ApJS, 74, 833
Hinkle, D. V. 1969, Biometrika, 56, 635
Hong, S., Calzetti, D., Dopita, M. A., et al. 2011, ApJ, 731, 45 (H11)
Hong, S., Katz, N., Davé, R., et al. 2010, arXiv:1008.4242v2
Hopkins, P. F., Quataert, E., & Murray, N. 2012, MNRAS, 421, 3522 (HQM)
Jarrett, T. H., Chester, T., Cutri, R., Schneider, S. E., & Huchra, J. P. 2003, AJ, 125, 525
Kauffmann, G., Heckman, T. M., Tremonti, C., et al. 2003, MNRAS, 346, 1055
Kennicutt, R. C. 1998, ARA&A, 36, 189
Kennicutt, R. C., Lee, J. C., Funes, S. J., et al. 2008, ApJS, 178, 247
Kewley, L. J., & Dopita, M. A. 2002, ApJS, 142, 35
Kewley, L. J., Dopita, M. A., Sutherland, R. S., Heisler, C. A., & Trevena, J. 2001, ApJ, 556, 121 (K01)
Kobulnicky, H. A., Kennicutt, R. C., Jr., & Pizagno, J. L. 1999, ApJ, 514, 544
Kobulnicky, H. A., & Skillman, E. D. 1996, ApJ, 471, 211
Kobulnicky, H. A., Skillman, E. D., Roy, J. R., Walsh, J. R., & Rosa, M. R. 1997, ApJ, 477, 679
Leitherer, C., Schaerer, D., Goldader, J. D., et al. 1999, ApJS, 123, 3
Lejeune, Th., Cuisinier, F., & Buser, R. 1997, A&A, 325, 229
MacLow, M.-M., & Ferrara, A. 1999, ApJ, 513, 142
MacLow, M.-M., McCray, R., & Norman, M. L. 1989, ApJ, 337, 141
Martin, C. L. 1997, ApJ, 491, 561
Martin, C. L., Kobulnicky, H. A., & Heckman, T. M. 2002, ApJ, 574, 663
Moustakas, J., Kennicutt, R. C., Jr., Tremonti, C. A., et al. 2010, ApJS, 190, 233
Murray, N., Menard, B., & Thompson, T. A. 2011, ApJ, 735, 66
Noeske, K. G., Weiner, B. J., Faber, S. M., et al. 2007, ApJL, 660, L43
Oppenheimer, B. D., & Davé, R. 2006, MNRAS, 373, 1265
Pilyugin, L. S., Vílchez, J. M., & Contini, T. 2004, A&A, 425, 849
Press, W. H., Teukolsky, S. A., Vetterling, W. T., & Flannery, B. P. 1992, Numerical Recipes The Art of Scientific Computing (2nd ed.; Cambridge: Cambridge Univ. Press)
Reynolds, R. J., Haftner, L. M., & Tuve, S. L. 1999, ApJL, 525, L21
Reynolds, R. J., Sterling, N. C., Haftner, L. M., & Tuve, S. L. 2001, ApJL, 548, L221
Rich, J. A., Dopita, M. A., Kewley, L. J., & Rupke, D. S. N. 2010, ApJL, 721, 505
Scannapieco, C., Tissera, P. B., White, S. D. M., & Springel, V. 2008, MNRAS, 389, 1137
Schaller, G., Schaerer, D., Meynet, G., & Maeder, A. 1992, A&AS, 96, 269
Silich, S., Tenorio-Tagle, G., & Munon-Tunon, C. 2003, ApJ, 590, 791
Soto, K. T., Martin, C. L., Prescott, M. K. M., & Armus, L. 2012, ApJL, 757, 86
Strickland, D. K., Heckman, T. M., Colbert, E. J. M., Hoopes, C. G., & Weaver, K. A. 2004, ApJL, 606, 829
Strickland, D. K., & Stevens, I. R. 2000, MNRAS, 314, 51
Suchkov, A. A., Berman, V. G., Heckman, T. M., & Balsara, D. S. 1996, ApJL, 463, 528
Sutherland, R. S., & Dopita, M. A. 1993, ApJS, 88, 253
Tremonti, C. A., Heckman, T. M., Kauffmann, G., et al. 2004, ApJL, 613, 898
Veilleux, S., & Osterbrock, D. E. 1987, ApJL, 63, 295
Weaver, R., McCray, R., & Castor, J. 1977, ApJ, 218, 377