On exact special solutions for the stochastic regularized long wave-Burgers equation

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Abstract
In this paper, we will analyze the Regularized Long Wave-Burgers equation with conformable derivative (cd). Some white noise functional solutions for this equation are obtained by using white noise analysis, Hermite transforms, and the modified sub-equation method. These solutions include exact stochastic trigonometric functions, hyperbolic functions solutions and wave solutions.

This study emphasizes that the modified fractional sub-equation method is sufficient to solve the stochastic nonlinear equations in mathematical physics.

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1 Introduction
Recently, fractional calculus gained considerable interests and significant theoretical developments in many fields and many studies have been achieved in this field [1–14]. Due to the fact that the stochastic models are more realistic than the deterministic models, we concentrate our study in this paper on the Wick-type stochastic time-fractional Regularized Long Wave-Burgers equation (RLWBE) with conformable derivative (cd). A lot of research on stochastic fractional differential equations has been done recently [15–18]. Ghany and Hyder [15] obtained analytical solutions to stochastic time-fractional KdV equations of the Wick-type, Ghany and Zakarya [16] obtained exact traveling wave solutions to a stochastic Schamel KdV equation of Wick-type, in [17] is analyzed a stochastic fractional KdV equation with cd, in [18] is used a white noise functional approach for the fractional coupled KdV equations and are obtained new soliton solutions. In this paper, we will analyze the time-fractional RLWBE.

The RLWBE with the aid of cd is given by [19, 20]

\[
D_\eta^\alpha p(x,\tau) + \delta D_\eta^\alpha p(x,\tau) + \epsilon(\tau)p(x,\tau)D_\eta^\alpha p(x,\tau) + \lambda(\tau)D_\eta^\alpha p(x,\tau) + \psi D_\eta^\alpha p(x,\tau) = 0, \\
t > 0, \quad 0 < \eta \leq 1, \\
(x, \tau) \in \mathbb{R} \times \mathbb{R}^+, \quad 0 < \eta \leq 1,
\]

(1.1)
where $\varepsilon(\tau), \lambda(\tau)$ and are limited measurable or integrable functions on $\mathbb{R}_+$, $D^\eta_x P(x, \tau)$ is the cd operator and $\delta$ and $\psi$ are real valued constants. In [19] are obtained some solutions of this equation by using the modified Kudryashov method. Zhao and Xuan [20] investigated the convergence and existence conditions of solutions for the RLWBE. Bona et al. [21] analyzed the integer ordered type of this equation for an evaluation modeled for water waves. Zhou and Liu [22] obtained kink type waves for the RLWBE. Inan et al. [23] acquired the hyperbolic and trigonometric solutions for this equation. Bas and Kilic [24] obtained the complex solutions by using an algebraic method.

The cd operator was exposed in [25]. This derivative operator can reform the failures of the other definitions. This important operator is the easiest, most natural and effectual definition of the fractional derivative for order $\eta \in (0, 1)$. We should note that the definition can be generalized to involve any $\eta$. All the same, the order $\eta \in (0, 1)$ is the most influential order.

We say that the conformable fractional differentiability of a function $f : [0, \infty) \mapsto \mathbb{R}$ is nothing else than the classical differentiability. Clearly, the conformable $\eta$-derivative of $f$ at some point $x > 0$, where $0 < \eta < 1$ is the pointwise product $x^{1-\eta}f'(x)$ [26].

The cd of order $\eta \in (0, 1)$ is described by the following statement [25]:

$$D^\eta f(t) = \lim_{\delta \to 0} \frac{f(t + \delta t^{1-\eta}) - f(t)}{\delta}, \quad f : (0, \infty) \to \mathbb{R}. $$

This derivative is more advantageous than others because it is easy to apply. Recently, there have been several researches on the conformable form of fractional calculations [27–31]. The stochastic model of Eq. (1.1) in the Wick sense with conformable derivatives can be given in the following process:

$$D^\eta x P(\delta + \varepsilon(\tau) \diamond P) \diamond D^\eta P + \lambda(\tau) \diamond D^2_{xx} P + \psi \diamond D^3_{xxt} P = 0, \quad (1.2)$$

where “$\diamond$” is the Wick product on the Kondratiev distribution space $(S)_{-1}$, $\varepsilon(\tau)$ and $\lambda(\tau)$ are $(S)_{-1}$-valued functions [18].

In order to obtain the exact solutions of the random RLWBE with conformable derivative, we only consider it in a white noise environment, that is, we will discuss the Wick-type stochastic RLWBE (1.2).

Our aim in this work is to obtain a new stochastic soliton and periodic wave solutions of the Wick-type stochastic RLWBE with the aid of cd. We use the modified sub-equation method [32, 33], white noise theory, and Hermite transform to produce a new set of exact soliton and periodic wave solutions for the RLWBE with cd. Moreover, we apply the inverse Hermite transform to obtain stochastic soliton and periodic wave solutions of the Wick-type stochastic RLWBE with the aid of cd. Finally, by an application example, we show how the stochastic solutions can be given as Brownian motion functional solutions.
2 Exact solutions of Eq. (1.1)

In this part, we will investigate exact solutions of RLWBE. Using the Hermite transform of Eq. (1.1), we use the deterministic equation

\[
D_\eta \tilde{P}(\kappa, \tau, z) + \tilde{\lambda}(\tau, z)D_{xx}^{\delta_{\delta_{\lambda}}} \tilde{P}(\kappa, \tau, z) + \tilde{\rho}(\tau, z)D_{xx}^{\delta_{\rho}} \tilde{P}(\kappa, \tau, z) = 0, \tag{2.1}
\]

where \(z = (z_1, z_2, \ldots) \in (\mathbb{C}^N)_c\) is a parameter. To obtain traveling wave solutions to Eq. (2.1), we introduce the transformations \(\tilde{\varepsilon}(\tau, z) = \varepsilon(\tau, z), \tilde{\lambda}(\tau, z) = \lambda(\tau, z), \tilde{P}(\kappa, \tau, z) = p(\kappa, \tau, z) = p(\xi(\kappa, \tau, z))\) with

\[
\xi(\kappa, \tau, z) = k \left(\frac{x^n}{\eta}\right) + \sigma \int_a^t \theta(\tau, z) d\tau
\]

where \(k, \sigma\) are arbitrary constants and \(\theta\) is a nonzero function to be determined. Hence, Eq. (2.1) can be converted to the following NODE:

\[
[\sigma \theta + k(\delta + \varepsilon(\tau, z)\rho)] \frac{dp}{d\xi} + \lambda(\tau, z)k^2 \frac{d^2p}{d\xi^2} + \rho k^2 \sigma \theta \frac{d^3p}{d\xi^3} = 0. \tag{2.2}
\]

- Considering the solution of Eq. (2.2), we can write it as a series expansion solution as follows:

\[
p(\xi) = \sum_{i=0}^{N} \alpha_i(\tau, z)G^i(\xi) + \sum_{i=1}^{N} \beta_i(\tau, z)G^{-i}(\xi), \tag{2.3}
\]

where \(\alpha_i\) \((i = 0, 1, \ldots, n)\), \(\beta_i\) \((i = 1, 2, \ldots, n)\) are functions to be determined later and \(G(\xi)\) satisfies the Riccati equation as follows:

\[
G'(\xi) = \sigma + G^2(\xi), \tag{2.4}
\]

where \(\sigma\) is an arbitrary constants.

- \(N\) is obtained with the aid of a balance between the highest order derivatives and the nonlinear terms in Eq. (2.2).

A few special solutions of Eq. (2.4) are given by [34]:

1. When \(\sigma < 0,\)

\[
G_1(\xi) = -\sqrt{-\sigma} \tanh(\sqrt{-\sigma} \xi), \quad G_2(\xi) = -\sqrt{-\sigma} \coth(\sqrt{-\sigma} \xi).
\]

2. When \(\sigma > 0,\)

\[
G_3(\xi) = \sqrt{\sigma} \tan(\sqrt{\sigma} \xi), \quad G_4(\xi) = \sqrt{\sigma} \cot(\sqrt{\sigma} \xi).
\]

3. When \(\sigma = 0, \rho = \text{const.},\)

\[
G_5(\xi) = -\frac{\Gamma(1 + \eta)}{\xi^n + \rho}. \tag{2.5}
\]
Remark. The generalized trigonometric and hyperbolic functions are defined as [2]

\[
\tan_\eta(\xi) = \frac{E_\eta(i\xi^n) - E_\eta(-i\xi^n)}{i(E_\eta(i\xi^n) + E_\eta(-i\xi^n))}, \quad \cot_\eta(\xi) = \frac{i(E_\eta(i\xi^n) + E_\eta(-i\xi^n))}{E_\eta(i\xi^n) - E_\eta(-i\xi^n)},
\]

\[
\tanh_\eta(\xi) = \frac{E_\eta(\xi^n) - E_\eta(-\xi^n)}{E_\eta(\xi^n) + E_\eta(-\xi^n)}, \quad \coth_\eta(\xi) = \frac{E_\eta(\xi^n) + E_\eta(-\xi^n)}{E_\eta(\xi^n) - E_\eta(-\xi^n)},
\]

where \(E_\eta(\xi) = \sum_{n=0}^{\infty} k^n \xi^n\) is the Mittag-Leffler function.

By balancing \(\frac{dp}{d\xi}\) with \(\frac{d^3p}{d\xi^3}\) in Eq. (2.2), it is found that \(N = 2\). Then we can choose the solution of Eq. (2.2) to be given by

\[
p(\xi) = a_0 + \alpha_1 G(\xi) + \alpha_2 G^2(\xi) + \beta_1 G^{-1}(\xi) + \beta_2 G^{-2}(\xi),
\]

where \(G(\xi)\) satisfies Eq. (2.4).

Now, replacing (2.7) and (2.4) into (2.2), by equating all coefficients of \(G(\xi)\), we can solve the equations. Then we obtain the following groups of solutions.

One of the obtained these groups is given by

\[
\alpha_0 = -\frac{\delta}{\varepsilon(\tau, z)} - \frac{12k^2\lambda(\tau, z)^2}{25\varepsilon(\tau, z)^2} \alpha_2 + \frac{1}{12} \left( \frac{1}{k^2\psi} + 8\sigma \right) \alpha_2,
\]
\[
\alpha_1 = -\frac{12k\lambda(\tau, z)}{5\varepsilon(\tau, z)}, \quad \beta_1 = \frac{4k\lambda(\tau, z)(-25\varepsilon(\tau, z)^2) - \frac{36k^2\lambda(\tau, z)^2}{\alpha_2^2}}{125\varepsilon(\tau, z)^3},
\]
\[
\beta_2 = \frac{864k^4\lambda(\tau, z)^4}{\varepsilon(\tau, z)^4} - \frac{300k^2\lambda(\tau, z)^2\sigma_2^2}{\varepsilon(\tau, z)^2} - \frac{625\sigma^2\alpha_2^4}{6875\alpha_2^2},
\]
\[
\sigma = -\frac{\varepsilon(\tau, z)\alpha_2}{12k\psi}.
\]

The exact solutions of Eq. (2.1) are given by:

(1) When \(\sigma < 0\),

\[
p_1(\kappa, \tau, z) = -\frac{\delta}{\varepsilon(\tau, z)} - \frac{12k^2\lambda(\tau, z)^2}{25\varepsilon(\tau, z)^2} \alpha_2 + \frac{1}{12} \left( \frac{1}{k^2\psi} + 8\sigma \right) \alpha_2
\]
\[
+ \frac{12k\lambda(\tau, z)}{5\varepsilon(\tau, z)} \sqrt{-\sigma} \tanh_\eta \left( \sqrt{-\sigma} \left( k \frac{\kappa^n}{\eta} \right) \right)
\]
\[
- \sigma \int_{\alpha}^{t} \frac{e^{\frac{\varepsilon(\tau, z)\lambda(\tau, z)}{12k\psi}}}{\eta^{1-\eta}} d\tau \right) \right)
\]
\[
\times \coth_\eta \left( \sqrt{-\sigma} \left( k \frac{\kappa^n}{\eta} \right) \right)
\]
\[
\times \frac{864k^4\lambda(\tau, z)^4}{\varepsilon(\tau, z)^4} - \frac{300k^2\lambda(\tau, z)^2\sigma_2^2}{\varepsilon(\tau, z)^2} - \frac{625\sigma^2\alpha_2^4}{6875\alpha_2^2} - \coth_\eta \left( \sqrt{-\sigma} \left( k \frac{\kappa^n}{\eta} \right) \right) \right),
\]
\[
p_2(\kappa, \tau, z) = -\frac{\delta}{\varepsilon(\tau, z)} - \frac{12k^2\lambda(\tau, z)^2}{25\varepsilon(\tau, z)^2} \alpha_2 + \frac{1}{12} \left( \frac{1}{k^2\psi} + 8\sigma \right) \alpha_2
\]
\[+ \frac{12k_2}{5e(\tau, z)} \sqrt{-\sigma} \coth(\sqrt{-\sigma} \left( k \left( \frac{x}{\eta} \right) - \sigma \int_a^t \frac{y(\tau, z, \eta)}{1-\eta} \, d\tau \right)) \]

\[- 4k \lambda(\tau, z)(-2\lambda(\tau, z)^2 - \frac{2\lambda(\tau, z)^2}{\sigma^2}) \frac{\tanh(\sqrt{-\sigma} \left( k \left( \frac{x}{\eta} \right) \right)}{125e(\tau, z)^3} \right) \]

\[- \sigma \int_a^t \frac{e^{(\tau, z, \eta)} y(t, z, \eta)}{1-\eta} \, d\tau \right) \]

\[\times \frac{364k_4(\tau, z)^4}{e(\tau, z)^4} - \frac{300k_4(\tau, z)^4 \sigma^2}{e(\tau, z)^2} - 625\sigma^2 \sigma^4 \]

\[\times \frac{1}{125e(\tau, z)^3} \right) \]

\[(2.10) \]

(2) When \( \sigma > 0 \),

\[p_2(\tau, z, \eta) = - \frac{\delta}{e(\tau, z)} - \frac{12k_2}{25e(\tau, z)^2} \alpha_2 + \left( \frac{1}{12} k^2 \psi + 8\sigma \right) \alpha_2 \]

\[- \frac{12k_2}{5e(\tau, z)} \sqrt{-\sigma} \tanh(\sqrt{-\sigma} \left( k \left( \frac{x}{\eta} \right) - \sigma \int_a^t \frac{y(\tau, z, \eta)}{1-\eta} \, d\tau \right)) \]

\[4k \lambda(\tau, z)(-2\lambda(\tau, z)^2 - \frac{2\lambda(\tau, z)^2}{\sigma^2}) \frac{\coth(\sqrt{-\sigma} \left( k \left( \frac{x}{\eta} \right) \right)}{125e(\tau, z)^3} \]

\[- \sigma \int_a^t \frac{e^{(\tau, z, \eta)} y(t, z, \eta)}{1-\eta} \, d\tau \right) \]

\[\times \frac{364k_4(\tau, z)^4}{e(\tau, z)^4} - \frac{300k_4(\tau, z)^4 \sigma^2}{e(\tau, z)^2} - 625\sigma^2 \sigma^4 \]

\[(2.11) \]

(3) When \( \sigma = 0, \rho = \text{const.} \),

\[p_5(\tau, z, \eta) = \left( \frac{1}{e(\tau, z)} - \frac{12k_2}{25e(\tau, z)^2} \alpha_2 + \left( \frac{1}{12} k^2 \psi + 8\sigma \right) \alpha_2 \right. \]

\[+ \frac{12k_2}{5e(\tau, z)} \frac{\Gamma(1 + \eta)}{((k \left( \frac{x}{\eta} \right) - \sigma \int_a^t \frac{y(\tau, z, \eta)}{1-\eta} \, d\tau)^n + \rho) \right)} \]

\[4k \lambda(\tau, z)(-2\lambda(\tau, z)^2 - \frac{2\lambda(\tau, z)^2}{\sigma^2}) \frac{\coth(\sqrt{-\sigma} \left( k \left( \frac{x}{\eta} \right) \right)}{125e(\tau, z)^3} \]

\[- \sigma \int_a^t \frac{e^{(\tau, z, \eta)} y(t, z, \eta)}{1-\eta} \, d\tau \right) \]

\[(2.12) \]
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\[ P_{1}(\kappa, \tau) = -\delta \frac{\varepsilon(\tau)}{\varepsilon(\tau)} - \frac{12k^{2}\lambda(\tau)\varepsilon(\tau)}{25e(\tau)\psi^{2}\varepsilon(\tau)} + \frac{1}{12} \left( \frac{1}{k^{2}\psi} + 8\sigma \right) \psi \lambda(\tau) \varepsilon(\tau) \]

\[ + \frac{12k^{2}\lambda(\tau)}{5e(\tau)} \sqrt{-\sigma} \tan h_{\eta} \left( \sqrt{-\sigma} \left( k \left( \frac{x^{\eta}}{\eta} \right) - \sigma \int_{a}^{t} \frac{\varepsilon(\tau)\psi(\tau)}{\eta^{2}\varepsilon(\tau)} d\tau \right) \right) \]

\[ - \frac{4k\lambda(\tau)\varepsilon(\tau)\psi(\tau)}{125e(\tau)} \sqrt{-\sigma} \coth_{\eta} \left( \sqrt{-\sigma} \left( k \left( \frac{x^{\eta}}{\eta} \right) \right) \right) \]

\[ - \sigma \int_{a}^{t} \frac{e(\tau)\psi(\tau)}{\eta^{2}\varepsilon(\tau)} d\tau \right) \]

\[ \times \coth_{\eta} \left( \sqrt{-\sigma} \left( k \left( \frac{x^{\eta}}{\eta} \right) \right) - \sigma \int_{a}^{t} \frac{e(\tau)\psi(\tau)}{\eta^{2}\varepsilon(\tau)} d\tau \right) \right) \]

\[ P_{2}(\kappa, \tau) = -\delta \frac{\varepsilon(\tau)}{\varepsilon(\tau)} - \frac{12k^{2}\lambda(\tau)\varepsilon(\tau)}{25e(\tau)\psi^{2}\varepsilon(\tau)} + \frac{1}{12} \left( \frac{1}{k^{2}\psi} + 8\sigma \right) \psi \lambda(\tau) \varepsilon(\tau) \]

\[ + \frac{12k^{2}\lambda(\tau)}{5e(\tau)} \sqrt{-\sigma} \coth_{\eta} \left( \sqrt{-\sigma} \left( k \left( \frac{x^{\eta}}{\eta} \right) - \sigma \int_{a}^{t} \frac{e(\tau)\psi(\tau)}{\eta^{2}\varepsilon(\tau)} d\tau \right) \right) \]

\[ - \frac{4k\lambda(\tau)\varepsilon(\tau)\psi(\tau)}{125e(\tau)} \sqrt{-\sigma} \tanh_{\eta} \left( \sqrt{-\sigma} \left( k \left( \frac{x^{\eta}}{\eta} \right) \right) \right) \]

3.333333333333333

\[ (\omega)^{2} \]

\[ (\lambda)^{2} \]

\[ (\sigma)^{2} \]
\[ -\sigma \int_{-\infty}^{t} \left( \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} - \frac{\phi(\tau)}{\tau^{1+\eta}} \right) d\tau \right) \]

\[ = \frac{-864k^4 \varepsilon(\tau)^{4\sigma}}{\varepsilon(\tau)^{4\sigma}} = \frac{300k^2 \varepsilon(\tau)^{2\sigma} \psi(\tau)^{2\sigma}}{\varepsilon(\tau)^{2\sigma}} - 625\sigma^2 \left( \frac{\psi(\tau)}{\varepsilon(\tau)} \right)^{2\sigma} \]

\[ \times \tan_1^2 \left( \sqrt{\frac{\sigma}{\varepsilon}} \left( \frac{\lambda(\tau)}{\eta} - \sigma \int_{a}^{t} \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} d\tau \right) \right). \tag{3.2} \]

(II) **Exact stochastic trigonometric solutions:**

\[ P_3(\varepsilon, \tau) = -\frac{\delta}{\varepsilon(\tau)} - \frac{12k^2 \varepsilon(\tau)^{2\sigma}}{25\varepsilon(\tau)^{2\sigma} \psi(\tau)^{2\sigma}} + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi(\tau)}{\varepsilon(\tau)} \]

\[ - \frac{12k \lambda(\tau)}{5\varepsilon(\tau)} \sqrt{\sigma} \tan_1 \left( \sqrt{\frac{\sigma}{\varepsilon}} \left( \frac{k}{\eta} \right) \right) \]

\[ - \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} \right) \right) + \frac{4k \lambda(\tau) \left( -25\varepsilon(\tau)^2 \sigma - \frac{36k^2 \varepsilon(\tau)^{2\sigma}}{\varepsilon(\tau)^{2\sigma}} \right)}{125\varepsilon(\tau)^{2\sigma}} \]

\[ \times \cot_1 \left( \sqrt{\frac{\sigma}{\varepsilon}} \left( \frac{k}{\eta} \right) - \sigma \int_{a}^{t} \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} d\tau \right) \right) \]

\[ + \frac{864k^4 \varepsilon(\tau)^{4\sigma}}{\varepsilon(\tau)^{4\sigma}} = \frac{300k^2 \varepsilon(\tau)^{2\sigma} \psi(\tau)^{2\sigma}}{\varepsilon(\tau)^{2\sigma}} - 625\sigma^2 \left( \frac{\psi(\tau)}{\varepsilon(\tau)} \right)^{2\sigma} \]

\[ \times \cot_1^2 \left( \sqrt{\frac{\sigma}{\varepsilon}} \left( \frac{k}{\eta} \right) - \sigma \int_{a}^{t} \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} d\tau \right) \right), \tag{3.3} \]

\[ P_4(\varepsilon, \tau) = -\frac{\delta}{\varepsilon(\tau)} - \frac{12k^2 \varepsilon(\tau)^{2\sigma}}{25\varepsilon(\tau)^{2\sigma} \psi(\tau)^{2\sigma}} + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi(\tau)}{\varepsilon(\tau)} \]

\[ - \frac{12k \lambda(\tau)}{5\varepsilon(\tau)} \sqrt{\sigma} \cot_1 \left( \sqrt{\frac{\sigma}{\varepsilon}} \left( \frac{k}{\eta} \right) - \sigma \int_{a}^{t} \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} d\tau \right) \right) \]

\[ 4k \lambda(\tau) \left( -25\varepsilon(\tau)^2 \sigma - \frac{36k^2 \varepsilon(\tau)^{2\sigma}}{\varepsilon(\tau)^{2\sigma}} \right) \]

\[ + \frac{125\varepsilon(\tau)^{2\sigma}}{125\varepsilon(\tau)^{2\sigma}} \]

\[ \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} \right) \right) \right) \right) \]

\[ + \frac{864k^4 \varepsilon(\tau)^{4\sigma}}{\varepsilon(\tau)^{4\sigma}} = \frac{300k^2 \varepsilon(\tau)^{2\sigma} \psi(\tau)^{2\sigma}}{\varepsilon(\tau)^{2\sigma}} - 625\sigma^2 \left( \frac{\psi(\tau)}{\varepsilon(\tau)} \right)^{2\sigma} \]

\[ \times \tan_1 \left( \sqrt{\frac{\sigma}{\varepsilon}} \left( \frac{k}{\eta} \right) - \sigma \int_{a}^{t} \frac{\varepsilon(\tau) \psi(\tau)}{12k \psi} d\tau \right) \right) \right), \tag{3.4} \]
(III) Exact stochastic wave solutions:

\[
P_3(x, \tau) = -\frac{\delta}{\varepsilon(\tau)} - \frac{12k^2 \lambda(\tau)^2}{25e(\tau)^2} \frac{1}{\varepsilon(\tau)} + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \left( \varphi_\lambda(\tau) \right) \frac{\psi_\lambda(\tau)}{\sigma \varepsilon(\tau)} + \frac{12k \lambda(\tau)}{5e(\tau)} \Gamma(1 + \eta) \frac{1}{\varepsilon(\tau)} \left( (k(\frac{x}{\eta}) - \sigma \int_a^t \frac{\psi_\lambda(\tau)}{\tau^{1-\eta}} d\tau)^\eta + \rho \right) + \frac{4k \lambda(\tau)}{125e(\tau)^2} \sigma \lambda(\tau) \frac{1}{\varepsilon(\tau)^2} \left( \frac{\psi_\lambda(\tau)}{\sigma \varepsilon(\tau)} \right)^2 \sigma \varepsilon(\tau) + \frac{864k^2 \lambda^2(\tau)^2 \sigma \lambda(\tau) \frac{1}{\varepsilon(\tau)^2} \left( \frac{\psi_\lambda(\tau)}{\sigma \varepsilon(\tau)} \right)^2}{6875 \frac{\psi_\lambda(\tau)}{\sigma \varepsilon(\tau)^2}} \Gamma(1 + \eta)^2 \eta \left( k(\frac{x}{\eta}) \right) \frac{\psi_\lambda(\tau)}{\sigma \varepsilon(\tau)} \frac{1}{12k^2 \psi} + \sigma \int_a^t \frac{\psi_\lambda(\tau)}{\tau^{1-\eta}} d\tau \right)^\eta + \rho \right)^2.
\]

(3.5)

4 Example

In this section, we investigate a special application example to represent the availability of our results and to confirm the real assistance of these results. We explain that the solutions of Eq. (1.2) are strongly dependent on the form of the given functions \(\varepsilon(\tau)\) and \(\lambda(\tau)\). So, for dissimilar forms of \(\varepsilon(\tau)\) and \(\lambda(\tau)\), we can find dissimilar solutions of Eq. (1.2) which come from Eqs. (3.1)–(3.5). We illustrate this by giving the following example.

When \(\eta = 1\),

\[
\tan_\eta(x) = \tan(x), \quad \cot_\eta(x) = \cot(x), \quad \tanh_\eta(x) = \tanh(x),
\]

\[
\coth_\eta(x) = \coth(x), \quad E_\eta(x) = \exp(x).
\]

Suppose \(\lambda(\tau) = \partial_x \varepsilon(\tau)\) and \(\varepsilon(\tau) = f(\tau) + \rho W_\tau\), where \(\partial\) and \(\rho\) are arbitrary constants, \(f(\tau)\) is a limited measurable function on \(\mathbb{R}_+\) and \(W_\tau\) is the Gaussian white noise which is the time derivative (in the strong sense in \((S, \gamma)\)) of the Brownian motion \(B_\tau\). The Hermite transform of \(W_\tau\) is given by \(\tilde{W}_\tau(z) = \sum_{n=0}^\infty \xi_n z^n \psi_n(\tau) d\tau [36]\). Using the definition of \(\tilde{W}_\tau(z)\), Eqs. (3.1)–(3.5) yield the white noise functional solution of Eq. (1.1) as follows:

\[
P_1(x, \tau) = -\frac{\delta}{f(\tau) + \rho W_\tau} - \frac{12k^2 \partial^2}{25f(\tau)^2} + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi_\lambda(\tau)}{\sigma} + \frac{12k \partial}{5} \sqrt{-\sigma} \tan \left( \sqrt{-\sigma}(kx - \frac{\partial}{12k} \left( \int_a^t f(\tau) d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) + c \right) \right)
\]

\[
+ \frac{4k \partial(-25\sigma - \frac{36k^2 \sigma^2}{\psi})}{125\sqrt{-\sigma}} \times \coth \left( \sqrt{-\sigma}(kx - \frac{\partial}{12k} \left( \int_a^t f(\tau) d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) + c \right) \right)
\]

\[
+ \frac{\partial}{12k} \left( \int_a^t f(\tau) d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) + c \right)
\]

\[
+ \frac{1}{12k} \left( \int_a^t f(\tau) d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) + c \right)
\]

\[
+ \frac{1}{12k} \left( \int_a^t f(\tau) d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) + c \right)
\]

\[
+ \frac{1}{12k} \left( \int_a^t f(\tau) d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) + c \right)
\]
\[- \frac{864k^4 \partial \sigma^3 - 300k^2 \partial \sigma \psi^2 \sigma - 625\sigma^2 \partial \psi^4}{6875\psi^3 \sigma} \times \coth^2 \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right), \tag{4.1} \]

\[ P_2(\kappa, \tau) = - \frac{\delta}{f(\tau) + \rho \mathcal{W}_\tau} \left( \frac{-12k^2 \partial^2}{25 \sqrt{\psi}} \right) + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi \partial}{\sigma} \]
\[ + \frac{12k\partial}{5} \sqrt{-\sigma} \coth \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right) \]
\[ - \frac{4k\partial(-25\sigma - \frac{36k^2 \sigma^2}{\psi})}{125 \sqrt{-\sigma}} \times \tan \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right), \tag{4.2} \]

\[ P_3(\kappa, \tau) = - \frac{\delta}{f(\tau) + \rho \mathcal{W}_\tau} \left( \frac{-12k^2 \partial^2}{25 \sqrt{\psi}} \right) + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi \partial}{\sigma} \]
\[ - \frac{12k\partial}{5} \sqrt{-\sigma} \tan \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right) \]
\[ + \frac{4k\partial(-25\sigma - \frac{36k^2 \sigma^2}{\psi})}{125 \sqrt{-\sigma}} \times \cot \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right), \tag{4.3} \]

\[ P_4(\kappa, \tau) = - \frac{\delta}{f(\tau) + \rho \mathcal{W}_\tau} \left( \frac{-12k^2 \partial^2}{25 \sqrt{\psi}} \right) + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi \partial}{\sigma} \]
\[ - \frac{12k\partial}{5} \sqrt{-\sigma} \cot \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right) \]
\[ + \frac{4k\partial(-25\sigma - \frac{36k^2 \sigma^2}{\psi})}{125 \sqrt{-\sigma}} \times \tan \left( \sqrt{-\sigma} \left( kx - \frac{\partial}{12k} \left\{ \int_a^t f(\tau) \, d\tau + \rho \left( B_\tau - \frac{\tau^2}{2} \right) \right\} + c \right) \right), \tag{4.4} \]
\[ P_5(x, \tau) = -\frac{\delta}{f(\tau) + \rho \mathcal{W}_\tau} - \frac{12k^2 \partial^2}{25(\psi \partial \psi)} + \frac{1}{12} \left( \frac{1}{k^2 \psi} + 8\sigma \right) \frac{\psi \partial}{\partial \sigma} \]
\[ + \frac{12k \partial}{5} \left( kx - \frac{1}{12} \left( \int_\tau^t f(\tau) \, d\tau + \rho \left( B_t - \frac{\tau^2}{2} \right) \right) + c + \rho \right) \]
\[ - \frac{4k \partial (25\sigma - \frac{6\partial^2 \psi^2}{\psi^3 })}{125} \left( \left( kx - \frac{1}{12} \left( \int_\tau^t f(\tau) \, d\tau + \rho \left( B_t - \frac{\tau^2}{2} \right) \right) + c + \rho \right) \right) \]
\[ - \frac{864k^4 \partial \psi}{6875} \left( kx - \frac{1}{12} \left( \int_\tau^t f(\tau) \, d\tau + \rho \left( B_t - \frac{\tau^2}{2} \right) \right) + c + \rho \right)^2, \quad (4.5) \]

where we have already used the following relations [33]:

\[ \tan \hat{\phi} (B_t) = \tan \left( B_t - \frac{\tau^2}{2} \right), \]
\[ \cot \hat{\phi} (B_t) = \cot \left( B_t - \frac{\tau^2}{2} \right), \]
\[ \tanh \hat{\phi} (B_t) = \tanh \left( B_t - \frac{\tau^2}{2} \right), \]
\[ \coth \hat{\phi} (B_t) = \coth \left( B_t - \frac{\tau^2}{2} \right). \]

5 Final remarks
We analyzed the RLWBE with cd for deterministic and stochastic forms. In addition, we studied a Wick-model stochastic RLWBE with cd. We investigate some exact solutions with the aid of the modified sub-equation method, Hermite transform and white noise theory. We obtained stochastic hyperbolic and trigonometric wave solutions via the inverse Hermite transform. Furthermore, we investigate an example, to show the stochastic solutions can be obtained as Brownian motion functional solutions. Besides, if \( \eta = 1 \), then the stochastic solutions (4.1)–(4.5) give a new set of stochastic solutions for the Wick-model stochastic RLWBE with integer derivatives.

This study emphasizes that the modified sub-equation method is sufficient to solve the stochastic nonlinear equations in mathematical physics. The applied method in this paper is a standard, direct, and computerized method, which lets us do confusing and boring algebraic calculations. It is shown that the process can be also applied to other nonlinear stochastic differential equations in mathematical physics.

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