Theoretical Study on the Origins of the Gap Bowing in Mg$_x$Zn$_{1-x}$O Alloys

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ABSTRACT

The full potential linear muffin-tin orbital (FP-LMTO) method was applied to study the structural and electronic properties of the compounds MgO, ZnO and their alloy Mg$_x$Zn$_{1-x}$O in the zincblende and NaCl structures. Results are obtained using the local density approximation (LDA), the ground-state properties like lattice constant and bulk modulus obtained agree very well with experimental and other theoretical calculations. The effect of composition on lattice constant from Vegard’s law and the bulk modulus was investigated. The microscopic origins of the gap bowing were explained by using the approach of Zunger and co-workers. It is concluded that the energy band gap bowing is primarily due to chemical charge-transfer effect. Contribution of volume deformation and structural relaxation to the gap bowing parameter is found to be very small.

Keywords: FP-LMTO Method; Structural and Electronic Properties; Mg$_x$Zn$_{1-x}$O Alloy; Bowing

1. Introduction

II-VI semiconductors have been of growing interest because of their wide band gap character and the potential applications for optoelectronic devices.

Currently, devices composed from the binary oxides of Mg, Zn, Cd and related alloys are generating considerable interest as they can provide, in principle, an accessible direct band-gap range from around 2.3 eV to 7.7 eV [1].

This makes them promising candidates even for deep ultraviolet (UV) lighting applications [2,3].

The B1 phase is found to be stable over all (Mg,Zn)O compositions [4], as expected from the preferences of the binary oxides. However, there exists the serious problem of phase separation due to the large dissimilarity of stable crystal structure, that is, wurtzite for ZnO and rocksalt for MgO.

Many ab initio calculations of the parent compounds, i.e., MgO and ZnO can be found in the literature [4,5]. To the best of our knowledge, there is no theoretical work on the structural and electronic properties of Mg$_x$Zn$_{1-x}$O alloy using FP-LMTO method.

Very recently, Fritsh et al. [6] have calculated the electronic properties of the rocksalt Mg$_x$Zn$_{1-x}$O alloy using empirical pseudopotential method, and Amrani et al. [7] using FP-LAPW method.

Many calculations using (FP-LMTO) have been done to characterize the structural, electronic and optical properties in four phases of compounds: RuC [8], and CuCl, CuBr, Cul [9] and of the cubic Be$_x$Zn$_{1-x}$Se alloy in zincblende structure [10].

Concerning, this method FP-LMTO, “all electron” it is meant that all electrons in the solid are considered in the calculation of electron density and total energy (as opposed, for instance, to a pseudo-potential method, where only the valence electrons are considered), and the term “full potential” implies that no approximation is made to the shape of the electron density or the electronic potential.

The present study focus mainly on the compositions dependence on the structural and electronic properties of Mg$_x$Zn$_{1-x}$O ternary alloy in the NaCl structure, using the full-potential linear muffin-tin orbital (FP-LMTO) method, within the local-density approximation (LDA) scheme, to determine a set of physical parameters of MgO, ZnO and their ternary alloy, namely the optimized lattice constant, bulk modulus, energy band gap and gap bowing.

A brief description of the computational details and methodology are given in Section 2. The most relevant results obtained for the structural and electronic properties for Mg$_x$Zn$_{1-x}$O in rocksalt phase are presented and discussed in Section 3. The conclusion is given in Section 4.
2. Method of Calculations

The calculations reported here were carried out using the \textit{ab initio} full-potential linear muffin-tin orbital (FP-LMTO) method \cite{11} as implemented in the Lmtart code \cite{12}.

The exchange-correlation potential was treated by the local density approximation (LDA) developed by Perdew and wang \cite{13}. This is an improved method compared to previous (LMTO) methods.

The FP-LMTO method treats muffin-tin spheres and interstitial regions on the same footing, leading to improvements in the precision of the eigenvalues. At the same time, the FP-LMTO method, in which the space is divided into interstitial regions (IR) and non overlapping muffin-tin spheres (MTS) surrounding the atomic sites, uses a more complete basis than its predecessors. In the IR regions, the basis functions are represented by Fourier series. Inside the MTS spheres, the basis functions are represented in terms of numerical solutions of the radial Schrödinger equation for the spherical part of the potential multiplied by spherical harmonics.

The charge density and the potential are represented inside the MTS by spherical harmonics up to $l_{\text{max}} = 6$. The integrals over the Brillouin zone are performed up to 35 special k-points for binary compounds and 27 special k-points for the alloys in irreducible Brillouin zone (IBZ), using the Blöchl’s modified tetrahedron method \cite{14}. The self-consistent calculations are considered to be converged when the total energy of the system is stable within $10^{-5}$ Ry. In order to avoid the overlap of atomic spheres the MTS radius for each atomic position is taken to be different for each case.

Both the plane waves cut-off are varied to ensure the total energy convergence. The values of the sphere radii (MTS), number of plane waves (NPLW), used in our calculation are summarized in Table 1.

3. Results and Discussions

3.1. Structural Properties

We first calculated structural properties of the binary compounds ZnO and MgO in the rocksalt (B1) phases.

Then, the alloy was simulated for compositions $x = 0.25$, $0.50$ and $0.75$ by applying special Quasi-random Sheme proposed by Zunger \textit{et al.} \cite{15}. This scheme has been applied to many semiconductors alloys \cite{16-22} successfully. We calculated the equilibrium structural parameters (lattice constant and bulk modulus) for the parent binary compounds MgO, ZnO and their ternary alloy.

As for the semiconductor ternary alloy in the type $B_xA_{1-x}C$ we have started our FP-LMTO calculations of the structural properties with B1 structure. We have chosen the basic cubic cell as the unit cell.

In the unit cell there are four C anions and three A and one B, two A and two B, and one A and three B cations respectively, for $x = 0.25$, $0.50$ and $0.75$. And the atomic positions for $\text{Mg}_x\text{Zn}_1-x\text{O}$ are given in Table 2.

We have assumed that the atoms are located at the ideal lattice sites in ordered positions.

The calculated lattice constant, bulk modulus and first-order pressure derivatives of bulk modulus (B'), for the ZnO, MgO and $\text{Mg}_x\text{Zn}_{1-x}\text{O}$ alloy are summarized in Table 3. For the considered structures, we perform the structural optimization by minimizing the total energy with respect to unit cell parameters.

Table 1. The plane wave number PW, energy cutoff (in Ry) and the muffin-tin radius (RMT) (in a.u) used in calculations for binary MgO, ZnO and their alloy in zincblende (B3) and rocksalt (B1) structures.

|         | PW  | RMT (a.u) Zn | Cutoff (Ry) |
|---------|-----|-------------|-------------|
| ZnO (B3) (B1) | 2974 | 2.4564 | 136.62 |
| MgO (B3) (B1) | 2974 | 2.1657 | 174.54612 |
| Mg$_{0.75}$Zn$_{0.25}$O | 44472 | 2.1585 | 151.24 |
| Mg$_{0.50}$Zn$_{0.50}$O | 44472 | 1.89 | 130.2546 |
| Mg$_{0.25}$Zn$_{0.75}$O | 44472 | 2.091 | 196.8642 |

Table 2. Atomic positions in the Mg$_x$Zn$_{1-x}$O alloy.

| Composition (x) | Atom | Atomic positions |
|-----------------|------|------------------|
| 0.25 Mg         | Zn   | (1/2,1/2,1/2)    |
|                 | O    | (0,0,0), (0,1/2,1/2), (1/2,0,1/2), (1/2,1/2,0) |
| 0.50 Mg         | Zn   | (1/2,0,0), (0,1/2,0) |
|                 | O    | (0,0,0), (0,1/2,1/2), (1/2,0,1/2), (1/2,1/2,0) |
| 0.75 Mg         | Zn   | (1/2,0,0), (0,1/2,0) |
|                 | O    | (0,0,0), (0,1/2,1/2), (1/2,0,1/2), (1/2,1/2,0) |
Table 3. Calculated lattice parameter aeq (Å), bulk modulus B (GPa) and first-order pressure derivatives of bulk modulus (B') for the ZnO, MgO and MgₓZn₁₋ₓO alloy. *Ref [26], †Ref [27], ‡Ref [28], §Ref [29], ¶Ref [30], ††Ref [31], ‡‡Ref [32], †††Ref [33], ‡‡‡Ref [7].

| x     | aeq (Å) | aeq (Exp) | otherwork | B (GPa) | B'     | B (Exp) | other work |
|-------|---------|-----------|-----------|---------|--------|---------|------------|
| ZnO   | 4.220   | 4.283⁶, 4.271⁷, 4.345⁸, 4.223⁹ | 198.02    | 4.23    | 202.5⁸ | 166.7³  |            |
| 0.25  | 4.208   |           |           | 4.316   | 190.03² | 4.29    | 161.4⁶     |
| 0.5   | 4.196   |           |           | 4.298   | 186.62⁴ | 4.39    | 157.2²     |
| 0.75  | 4.185   |           |           | 4.279   | 179.08  | 4.36    | 153.2³     |
| MgO   | 4.173   | 4.213⁶, 4.212², 4.259⁶, 4.247⁸ | 175.06    | 4.02    | 156⁴   | 161.9⁹   |

The calculated lattice constant, bulk modulus and its pressure derivative for each x of the MgₓZn₁₋ₓO with respect to the cell parameters and also the atomic positions. The total energies were calculated as a function of volume and were fitted to the universal Murnaghan’s equation of state [23]. The predicted lattice parameters for the binary compounds are in reasonable agreement with those measured experimentally as well as with those calculated from theoretical methods.

However, we have a small underestimation of the lattice parameters, when we compare our results to the experimental data, this is due essentially to the use of the LDA. The calculated values of the bulk modulus decrease from ZnO to MgO, i.e. from the higher to the lower atomic number. This suggests that ZnO is more rigid than MgO.

Usually, in the treatment of alloy problems, it is assumed that the atoms are located at ideal lattice sites and the lattice constants of alloys should vary linearly with compositions x according to the so-called Vegard’s law [24]:

\[
a \left( A_x B_{1-x} C \right) = x a_{AC} + (1-x) a_{BC}
\]

where \(a_{AC}\) and \(a_{BC}\) are the equilibrium lattice constants of the binary compounds \(AC\) and \(BC\) respectively, and \(a(A_x B_{1-x} C)\) is the alloy lattice constant.

The calculated lattice constants at different compositions, as shown in Figure 1, where well fitted with the following relation:

\[
\left( A_x B_{1-x} C \right) = x a_{AC} + (1-x) a_{BC} - x(1-x)b
\]

de the quadratic term \(b\) represents the disorder parameter (bowing), the obtained bowing parameter \(b\) is 0.00114.

The linear dependence in \(x\) is therefore in accordance with Vegard’s Law with negligible bowing. Figure 2, show the variation of the bulk modulus versus concentration \(x\) for MgₓZn₁₋ₓO alloy.

The overall behaviors of the variation of the bulk modulus as a function of the composition for the MgₓZn₁₋ₓO is presented in Figure 2 is compared to the results predicted by linear concentration dependence (LCD). A deviation from the LCD is observed with downward bowing equal to 4.19 GPa is obtained by fitting a calculated data with a polynomial of second order. This deviation from Vegard Law [24] is due to the fact that bulk modulus of ZnO is 11.2% higher than that of MgO.

To the best of our knowledge, there are no experimental work exploring the structural properties (e.g, the bulk modulus B and it pressure derivatives B’) and the bowing parameters of the of the investigated alloys, but our results are relatively close to those of Ref. [7] obtained by using the full potential-linear augmented plane wave (FP-LAPW) method.
3.2. Electronic Properties

We have calculated the band structures for the Mg\textsubscript{x}Zn\textsubscript{1-x}O alloy along the high directions in the first Brillouin zone at the calculated equilibrium lattice constants.

The band structure calculations give a direct band gap  \( \Gamma - \Gamma \) for MgO and an indirect band gap  \( M - \Gamma \) for ZnO and Mg\textsubscript{x}Zn\textsubscript{1-x}O. The calculated band gaps for all studied compositions (\( x = 0, 0.25, 0.50, 0.75, 1 \)) are given in Table 4. We calculated the gap bowing by fitting the nonlinear variation of the calculated band gap versus composition \( x \) with the quadratic semi-empirical formula:

\[
E_g^{AC} = xE_g^{AC} + (1-x)E_g^{BC} - x(1-x)b_g
\]

where  \( E_g^{AC} \) and  \( E_g^{BC} \) are the energy band gaps of the ternary alloy  \( A_xB_{1-x}C \) and its binary parents  \( AC \) and  \( BC \), respectively. The curvature  \( b_g \) is commonly known as gap bowing parameter. The results shown in Figure 3, are well fitted by the expression (1), and are summarized as follows:

\[
\begin{align*}
\text{Mg}_x\text{Zn}_{1-x}\text{O}(M - \Gamma) & \rightarrow E_g(x) = 1.376 - 1.379x + 4.981x^2 \quad (2) \\
\text{Mg}_x\text{Zn}_{1-x}\text{O}(\Gamma - \Gamma) & \rightarrow E_g(x) = 2.754 + 0.104x + 1.861x^2 \quad (3)
\end{align*}
\]

It is clear from the above equations that the direct (\( \Gamma - \Gamma \)) and indirect (\( M - \Gamma \)) bands gaps versus concentration have a nonlinear behavior. This behavior was observed by Amrani et al. [7] by using  \textit{ab initio} FP-LAPW. The indirect gap has a downward bowing with a value of 4.98, and 1.86 for the direct gap.

From the Figure 3, it is clear that the crossover of indirect band gap (\( M - \Gamma \)) to direct (\( \Gamma - \Gamma \)) is at concentration of 0.9.

In order to better understand the physical origins of the gap bowing parameter in Mg\textsubscript{x}Zn\textsubscript{1-x}O alloys, we follow the procedure of Bernard and Zunger [25] and decompose the total bowing parameter  \( b \) into three physically distinct contributions. The overall bowing coefficient at each composition \( x \) measures the change in the band gap according to the formal reaction:

\[
x\text{MgO}(a_{\text{MgO}}) + (1-x)\text{ZnO}(a_{\text{ZnO}}) \rightarrow \text{Mg}_x\text{Zn}_{1-x}\text{O}(a_{eq})(4)
\]

where  \( a_{\text{MgO}} \) and  \( a_{\text{ZnO}} \) are the equilibrium lattice constants of the binary compounds and  \( a_{eq} \) is the equilibrium lattice constant of the alloy with average composition \( x \). Equation (4) is decomposed into three steps:

\[
\begin{align*}
\text{MgO}(a_{\text{MgO}}) + \text{ZnO}(a_{\text{ZnO}}) & \rightarrow \text{MgO}(a) + \text{ZnO}(a) \quad (5) \\
x\text{MgO}(a) + (1-x)\text{ZnO}(a) & \rightarrow \text{Mg}_x\text{Zn}_{1-x}\text{O}(a) \quad (6) \\
\text{Mg}_x\text{Zn}_{1-x}\text{O}(a) & \rightarrow \text{Mg}_x\text{Zn}_{1-x}\text{O}(a_{eq}) \quad (7)
\end{align*}
\]

The first step measures the volume deformation (VD) effect on the bowing. The corresponding contribution  \( b_{VD} \) to the bowing parameter represents the relative response of the band structure of the binary compounds MgO and ZnO to hydrostatic pressure, which here arises from the change of their individual equilibrium lattice constants to the alloy value  \( a = a(x) \). The second contribution, the charge exchange (CE) contribution  \( b_{CE} \), reflects the charge transfer effect which is due to the different (averaged) bonding behavior at the lattice constant  \( a \). The last contribution, the so called: structural relaxation (SR), measures changes in passing from the relaxed alloy by  \( b_{SR} \). Consequently, the total bowing parameter is defined as:

\[
b = b_{VD} + b_{CE} + b_{SR} \quad (8)
\]

The general representation of the composition-dependent band gap of the alloy in terms of the gaps of the binary compounds  \( \text{MgO} \) and  \( \text{ZnO} \) to hydrostatic pressure, which here arises from the change of their individual equilibrium lattice constants to the alloy value  \( a = a(x) \). The second contribution, the charge exchange (CE) contribution  \( b_{CE} \), reflects the charge transfer effect which is due to the different (averaged) bonding behavior at the lattice constant  \( a \). The last contribution, the so called: structural relaxation (SR), measures changes in passing from the relaxed alloy by  \( b_{SR} \). Consequently, the total bowing parameter is defined as:

\[
E_g(x) = xE_{\text{MgO}}(a_{\text{MgO}}) + (1-x)E_{\text{ZnO}}(a_{\text{ZnO}}) - bx(1-x) \quad (9)
\]

This allows a splitting of the total bowing parameter  \( b \) into three contributions according to:

\[
\begin{align*}
b_{VD} &= \frac{E_{\text{MgO}}(a_{\text{MgO}}) - E_{\text{MgO}}(a)}{1-x} + \frac{E_{\text{ZnO}}(a_{\text{ZnO}}) - E_{\text{ZnO}}(a)}{x} \quad (10) \\
b_{CE} &= \frac{E_{\text{MgO}}(a)}{1-x} + \frac{E_{\text{ZnO}}(a)}{x} - \frac{E_{\text{MgZnO}}(a)}{x(1-x)} \quad (11) \\
b_{SR} &= \frac{E_{\text{MgZnO}}(a) - E_{\text{MgZnO}}(a_{eq})}{x(1-x)} \quad (12)
\end{align*}
\]

where  \( E \) is the energy band gap calculated for the indicated compound with the indicated atomic positions and lattice constant  \( a_{\text{MgO}}, a_{\text{ZnO}} \) and  \( a_{eq} \) are the equilibrium lattice constants of MgO, ZnO and Mg\textsubscript{x}Zn\textsubscript{1-x}O alloys.
respectively. The lattice constant \( (a) \) is calculated by linear composition dependence rule \([24]\) for the alloys.

Using Equations (10)-(12), the bowing coefficients \( b \) calculated at molar fractions \( x = 0.25, 0.50 \) and 0.75 for the \( \text{Mg}_x \text{Zn}_{1-x} \text{O} \) alloy are listed in \( \text{Table 5} \), together with the bowing obtained using a quadratic variation of the band gap energy versus composition \( x \). The calculated quadratic parameters of the gap bowing are in good agreement with the values found from the approach of Bernard and Zunger \([25]\). The charge transfer contribution \( b_{\text{CE}} \) dominates the total gap bowing parameter in the three compositions \( x: (x = 0.25, 0.50 \) and 0.75); this is related to electronegativity mismatch between the constituting atoms: \( \text{Mg} (1.31), \text{Zn} (1.65) \) and \( \text{O} (3.44) \). The low value of \( b_{\text{VD}} \) is related to the weak mismatch of the lattice parameters of \( \text{MgO} \) and \( \text{ZnO} \) compounds. The small contribution of the structural relaxation to the bowing parameter it due to that our calculations are for ordered structure.

### 4. Conclusions

In this study, we have presented a complete theoretical analysis of the structural and electronic properties of \( \text{Mg}_x \text{Zn}_{1-x} \text{O} \) alloys by using the FP-LMTO method within the local-density approximation (LDA). The equilibrium lattice constants, bulk modulus and first order pressure derivatives of the bulk modulus of the binary compounds and \( \text{Mg}_x \text{Zn}_{1-x} \text{O} \) alloy have been studied.

The energy gaps of \( \text{ZnO} \) and \( \text{MgO} \) compounds calculated with the equilibrium lattice constant, are found to be reasonable agreement with the experimental data. We have investigated the composition dependence of the lattice constant, bulk modulus and band gap. The calculated lattice constants scale linearly with composition \( x \): this is related to electronegativity mismatch between the constituents: \( \text{Mg} (1.31), \text{Zn} (1.65) \) and \( \text{O} (3.44) \). The small contribution of the structural relaxation to the bowing parameter it due to that our calculations are for ordered structure.

| X         | \( \text{Γ - Γ} \) | \( \text{M - Γ} \) |
|-----------|---------------------|---------------------|
| \( \text{ZnO} \) | 2.715               | 2.60               | 2.55               | 1.96               | 2.45               | 1.1                 |
| \( 0.25 \)  | 2.956               | 2.90               | 1.644              | 1.74               |
| \( 0.50 \)  | 3.297               | 3.31               | 2.069              | 2.20               |
| \( 0.75 \)  | 3.780               | 3.89               | 2.581              | 2.88               |
| \( \text{MgO} \) | 4.766               | 7.8               | 4.98               | 5.40               | 5.130              | 5.98               |

### Table 5. Decomposition of optical bowing into volume deformation (VD), charge exchange (CE) and structural relaxation (SR). Contribution compared with the optical bowing obtained by a quadratic interpolation. (All values are in eV).

|                | Zunger approach | Quadratic equation | other work Ref. [7] |
|----------------|-----------------|--------------------|---------------------|
| \( \text{Mg}_{0.25} \text{Zn}_{0.75} \text{O} \) | bvp 0.415, bce 4.391, bsr 0.160 | bvp 0.311, bce 5.130, bsr −0.010, b 5.431, 4.98 | 6.206, 6.219 |
|                | \( \text{Mg}_{0.625} \text{Zn}_{0.375} \text{O} \) | bvp 0.519, bce 5.129, bsr −0.021, b 5.627, 4.98 |
|                | \( \text{Mg}_{0.625} \text{Zn}_{0.375} \text{O} \) | bvp 0.519, bce 5.129, bsr −0.021, b 5.627, 4.98 |

The main advantage with a full-potential linear muffin-tin orbital (FP-LMTO) method, as described here, is that the electronic structure problem is solved with very high accuracy, so that total energies can be calculated with high precision.

### Table 4. Indirect and direct band gap energy of \( \text{Mg}_x \text{Zn}_{1-x} \text{O} \) alloy. \( ^a \text{Ref}[34] \), \( ^b \text{Ref}[33] \), \( ^c \text{Ref}[34] \), \( ^d \text{Ref}[35] \), \( ^e \text{Ref}[36] \), \( ^f \text{Ref}[7] \).

| X         | \( \text{Γ - Γ} \) | \( \text{M - Γ} \) |
|-----------|---------------------|---------------------|
| \( \text{ZnO} \) | 2.715               | 2.60               | 2.55               | 1.96               | 2.45               | 1.1                 |
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| \( 0.50 \)  | 3.297               | 3.31               | 2.069              | 2.20               |
| \( 0.75 \)  | 3.780               | 3.89               | 2.581              | 2.88               |
| \( \text{MgO} \) | 4.766               | 7.8               | 4.98               | 5.40               | 5.130              | 5.98               |
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