The 2008-2009 outburst of the young binary system Z CMa unraveled by interferometry with high spectral resolution

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**ABSTRACT**

Z CMa is a young binary system consisting of an Herbig primary and a FU Ori companion. Both components seem to be surrounded by active accretion disks and a jet was associated to the Herbig B0. In Nov. 2008, K. Grankin discovered that Z CMa was exhibiting an outburst with an amplitude larger than any photometric variations recorded in the last 25 years. To study the innermost regions in which the outburst occurs and understand its origin, we have observed both binary components with AMBER/VLTI across the Br\textsubscript{γ} emission line in Dec. 2009 in medium and high spectral resolution modes. Our observations show that the Herbig Be, responsible for the increase of luminosity, also produces a strong Br\textsubscript{γ} emission, and they allow us to disentangle from various origins by locating the emission at each velocities through the line. Considering a model of a Keplerian disk alone fails at reproducing the asymmetric spectro-astrometric measurements, suggesting a major contribution from an outflow.

**Keywords:** Astrophysics, Young star, Circumstellar Matter, Interferometry, Infrared, Outburst

1. **INTRODUCTION**

Accretion plays an important role in star and planet formation. For many years, it was considered to be a slow quasi-stationary process\textsuperscript{1} occurring mostly through a viscous disk ending in its inner part by a boundary layer\textsuperscript{2} with the star or by magnetospheric funnels\textsuperscript{3,4}. However, this scenario has been challenged by observations\textsuperscript{5,6} that suggest that the accretion process could be time-variable and occur quickly by means of short high mass accretion rate bursts. Studying the very inner region of a young stellar object that is known to experience episodic photometric outbursts is thus of prime importance to understand the role of accretion in the formation of the star and its environment.

The study presented in this paper is part of a large observational campaign targeting Z CMa during its 2008 outburst that aims to understand its origin. To directly probe the morphology of the hot gas in the inner AU\textsubscript{s}, we took advantage of the spatial and spectral resolution available at the VLTI to perform micro-arcsecond spectro-astrometry. With AMBER, we resolve the K-band emission of the hot gas surrounding each star at the milliarcsecond resolution. This paper outlines the analysis already published by Benisty et al. (2010)\textsuperscript{7}.

Section 2 summarizes what we know about Z CMa. In Sect. 3, we present the observations and the data processing. Section 4 presents the findings of our study that we discuss in Sect. 5.

2. **THE SYSTEM Z CANIS MAJORIS**

Z CMa is a pre-main-sequence binary with a separation\textsuperscript{9,10} of 0.1" located at a distance estimated\textsuperscript{11,12} from 930 to 1150 pc. The primary, embedded in a dust cocoon, was identified as a Herbig Be star based on spectropolarimetry\textsuperscript{13}. It is surrounded by an inclined disk, possibly a circumbinary disk, as inferred from millimeter observations\textsuperscript{14} and dominates the infrared continuum and total luminosity of the system. In contrast, the secondary is the major source of continuum emission at visual wavelengths. Although the secondary has not
undergone a large outburst this century, it was identified as a FU Or object based on its broad double-peaked optical absorption lines, which are typical of a circumstellar disk that undergoes a strong accretion, and spectral type\textsuperscript{15} of F-G. In the past twenty years, the Z CMa system exhibited repeated brightness variations, of \~0.5-1 visual magnitude, which were attributed to the Herbig Be star\textsuperscript{16}. Z CMa is clearly associated with a bipolar outflow\textsuperscript{17,18} that extends to 3.6 pc along PA\~240°. A 1\"x0.24\" micro-jet\textsuperscript{18} was detected in the [OI] 6300 Å line in the same direction, and the authors of this work\textsuperscript{18} concluded that the optical emission-line spectrum and the jet are associated with the primary. However, the innermost environments of the Z CMa components have been poorly studied. Two broad-band interferometric measurements have been obtained, allowing only characteristic sizes\textsuperscript{20,21} of the K-band continuum emission to be derived.

In January 2008, Z CMa’s brightness increased\textsuperscript{8} by about two visual magnitudes, representing the largest outburst observed in the past 90 years (see Fig. 1). Based on spectropolarimetric observations\textsuperscript{22} this outburst was considered to be associated with the Herbig Be star. The overall spectral energy distribution of the system is strongly modified during the outburst at wavelengths shorter than 10\µm, which indicates that the outburst originates close to the star.

Figure 1. Light curve of Z CMa showing the 2008 outburst (see Grankin et al. 2009\textsuperscript{8}) complemented with the last data. Top panel: all data observed since 1930. Middle panel: the last 25 years. Bottom panel: the last ten years showing the sudden raise of \( V \) by two magnitudes in 2008 and which disappears beginning of 2010. Black filled circles: photometric observations from CrAO; green squares: observations from Mt. Maidenak; red crosses: CCD observations from ASAS; and blue points: visual estimations of a brightness from AAVSO.
Table 1. Log of the observations.

| Date     | Baseline | Projected length (m) | Position angle (°) | $R$   |
|----------|----------|----------------------|--------------------|-------|
| 05/12/08 | D0-G1    | 69                   | 137                | 1500  |
|          |          |                      |                    | FUOr+HBe |
| 07/12/08 | K0-G1    | 89                   | 28                 | 12000 |
|          |          |                      |                    | FUOr+HBe |
| 09/12/08 | K0-G1    | 88                   | 24                 | 12000 |
|          |          |                      |                    | FUOr+HBe |
| 15/12/08 | U2-U3    | 44                   | 35                 | 1500  |
|          | U3-U4    | 62                   | 107                |       |
|          | U2-U4    | 86                   | 78                 |       |
| 16/12/08 | U1-U2    | 56                   | 35                 | 1500  |
|          | U2-U4    | 77                   | 88                 |       |
|          | U1-U4    | 120                  | 66                 |       |
| 10/01/10 | D0-G1    | 71                   | 133                | 1500  |
|          |          |                      |                    | FUOr+HBe |

Note: $R$ is the spectral resolution. ‘FUOr+HBe’ specifies when the binary is in the field of view.

3. OBSERVATIONS AND DATA PROCESSING

Z CMa was observed at the Very Large Telescope Interferometer (VLTI), using the AMBER instrument that allows the simultaneous combination of three beams in the near-infrared. The instrument delivers spectrally dispersed interferometric observables (visibilities, closure phases, differential phases) at spectral resolutions up to 12,000.

We report $K$-band observations taken in the medium spectral resolution mode ($R \sim 1500$) with the 8.2 m Unit Telescopes (UTs) as well as with the 1.8 m Auxiliary Telescopes (ATs), and in the high spectral resolution mode ($R \sim 12,000$) with the ATs. The data were obtained within programs of Guaranteed Time, Director’s Discretionary Time, and Open Time observations. Z CMa was observed with 11 different baselines of 4 VLTI configurations, during 5 nights in December 2008 and one night in January 2010 (see Table 1 for the summary of observations; the baselines G1-A0 and K0-A0 which could not be used are not reported). The longest baseline is $\sim 120$ m corresponding to a maximum angular resolution of 3.7 mas. With the UTs, the observations are coupled with the use of adaptive optics and the resulting field of view ranges from 50 to 60 mas. This allowed us to spatially resolve the binary and obtain separate measurements of the FU Or and the Herbig Be. In contrast, the ATs field-of-view, ranging from 230 to 280 mas, includes both stars and the interferometric signal results from both emissions. In addition to Z CMa, calibrators (HD45420, HD60742, HD55137, HD55832) were observed to correct for instrumental effects. All observations were performed using the fringe-tracker FINITO.

The data reduction was performed following the standard procedures, using the amdlib package, release 2.99, and the yorick interface provided by the Jean-Marie Mariotti Center. Raw spectral visibilities, differential phases, and closure phases were extracted for all the frames of each observing file. A selection of 80% of the highest quality frames was made and to avoid the effects of instrumental jitter and unsatisfactory light injection. Consecutive observations were merged to enhance the signal-to-noise ratio. Calibration of the AMBER+VLTI instrumental transfer function was done using measurements of the calibrators, after correcting for their diameter. The accuracy of the wavelength/velocity calibration is $\sim 50$ km/s. Because the $K$-band continuum measured by the ATs (due to both stars) is very resolved on long baselines ($V^2 \sim 0$), the observations obtained on the G1-A0 and K0-A0 baselines could unfortunately not be exploited. The absolute value of the visibilities obtained with the UT baselines could not be determined due to random vibrations of the telescopes. However, this issue affects all spectral channels in the same way, and does not modify our conclusions.

4. FINDINGS

We recall that the visibilities provide information about the spatial extent of the emission, and decrease as the extension increases. Differential phases provide a measurement of the photocenter displacements across the sky, projected along the baseline direction. They can therefore be converted into differential spectro-astrometric

\*http://www.jmmc.fr/amberdrs
shifts. They are measured relative to the continuum, for which we assume a zero phase. Finally, the closure phases are related to the asymmetry of the brightness distribution (e.g., they are null for a point-symmetric object).

4.1 The Brγ line is present only during the outburst

Fig. 2 compares the spectra and the visibilities obtained during and after the outburst: the emission line, and the signature in the visibilities, disappear after the outburst. Plotted within a large velocity range, the visibilities show a typical signature of binarity (i.e., a cosine modulation), in agreement with the system main characteristics (separation, position angle, flux ratio; Bonnefoy et al., in prep.).

The main consequence of the outburst at least for interferometric observations is to produce a Brγ in the spectrum and a raise of the visibility which disappears when the outburst is over. Therefore, the origin of the outburst must be searched in the emission line, the Brγ line being well-suited.
Figure 3. Interferometric measurements of the Z CMa binary system with the UTs during the outburst in December 2008. Medium resolution spectra, spectrally dispersed squared visibilities normalized to the continuum ones $V_{c}^{2}$, differential phases ($\Delta \Phi$), and closure phases (CP) measured for the FU Or (left) and the same quantities for the Herbig Be (right). The binary system as observed by the NACO instrument at the VLT (adaptive optics with infrared camera) is displayed at the bottom with identification of the sources. The square boxes represent the field of view of these individual measurements.

4.2 The outburst is seen only in the Herbig Be component

The left and right columns of Fig. 3 present examples of the MR observations obtained with the UTs for each star during the outburst. Each column includes a spectrum (normalized to the continuum), squared visibilities, differential phases, and closure phases. Since the absolute values of the visibilities measured with the UTs are unknown, we normalized the continuum values to 1 – even though the emission is resolved.

For the FU Or (left panels), within the error bars, the spectrum shows Br$_{\gamma}$ in neither emission nor absorption. Consequently, no change in the visibilities or phases across the line is expected nor seen. In contrast, the Herbig Be star exhibits a clear Br$_{\gamma}$ line in emission (right panels), although at this spectral resolution ($\Delta v \sim 200$ km/s), the line is not spectrally resolved. The visibility increases through the line and the differential phases produce an S-shape variation. The closure phases differ from zero, with values of $25^\circ \pm 12^\circ$. The phase, line, and visibility signals are present from $\sim 600$ to $500$ km/s, although because of the low line-to-continuum ratio in the extended...
4.3 Size of the Herbig Be system in the continuum and in Br$_\gamma$

Figure 4 shows measurements obtained with the ATs, i.e., with both stars in the field of view. In this case, the level of continuum is determined by both stellar components. These panels present observations obtained in high spectral resolution ($R \sim 12000$). In this case, the line is spatially and spectrally resolved ($\Delta v \sim 25$ km/s), and the spectra exhibit a clear double-peaked and asymmetric profiles, with less emission at blueshifted velocities. The spectral visibilities present a similar profile.

From the visibilities, one can locate the emission at each velocity and distinguish between various scenarios capable of producing the line. The visibility increase within the line implies that the Br$_\gamma$ emitting region is more compact than the one responsible for the continuum. To derive the characteristic sizes of the region emitting Br$_\gamma$,
Figure 5. Differential phases measured with the UTs (black crosses and lines) for six different baselines. The dots in different colors represent various velocity channels, from dark blue (for \( v \sim -600 \) km/s) to dark red (for \( v \sim +500 \) km/s). The red line is the 2D astrometric solution \( \vec{p}(\lambda) \).

only, for each spectral channel of the HR measurements, one has to subtract the underlying continuum to first determine the visibility of the line only. These estimates can only be performed using the data gathered with the ATs, for which reliable absolute values for the Br\(_{\gamma}\) visibilities are obtained. Using a model of an uniform ring, the emission in the line has a typical extension (ring diameter) of \( \sim 1.6 \) mas at zero velocity, and \( \sim 2.5 \) mas at higher velocities (\( \sim 100 \) km/s), i.e., from \( \sim 1.5 \) to \( \sim 2.6 \) AU, depending on the distance.

As the continuum emission measured with the ATs includes both stars, it is not direct to establish the typical size of the Herbig Be continuum. In contrast, the UTs data includes only one stellar component. Although no absolute visibility values can be obtained, size ratios between the line and the continuum can be derived. Using the sizes previously estimated for the line from the ATs data, typical sizes of \( \sim 3.4 \) mas (\( \sim 3.6 \) AU) for the Herbig Be K-band continuum can be determined, in agreement with the previous estimate (\( \sim 3.9 \) mas in 2004).

Considering a dust sublimation temperature around 1500-2000 K and the stellar properties determined by van den Ancker et al. (2004), the inner edge of the dusty disk must be located at \( \sim 4-7 \) AU, in agreement with our findings. An asymmetry in the inclined inner disk could explain the non-zero closure phases measured at a level similar to other Herbig AeBe stars. Such methodology is however uncertain as it depends on the stellar contribution of the Herbig Be to the K-band continuum, that is unknown.

The Br\(_{\gamma}\) emitting region in the Herbig Be component is therefore more compact than the one responsible for the continuum.

4.4 Spectro-astrometry of the Br\(_{\gamma}\) line

The differential phases \( \Delta \Phi \) which are represented in Fig. can be expressed in terms of photocenter displacements \( p \) (in arcseconds), given by \( p = -2\pi \Delta \Phi / B \lambda \), where \( \lambda \) and \( B \) are the wavelength and the projected
baseline length of the observations, respectively. $p$ is the projection along the baseline direction, of the 2D photocenter vector $\vec{p}$ in the plane of the sky (i.e., of a spectro-astrometric signal).

We fitted all the differential phases along the 6 available baselines with a single vector $\vec{p}$, independently of each spectral channel. Figure 5 shows the differential phases and the best solution for $\vec{p}$. The left plot of Fig. 6 gives $\vec{p}$ in a 2D map of the plane of the sky. Clear asymmetric displacements, up to $\sim 150$ micro-arcseconds, are observed, both at red-shifted and blue-shifted velocities. In this case, $\vec{p}$ accounts for the emission of both the line and the continuum.

Subtracting the continuum contribution to determine the photocenter displacements, $\vec{p}_{Br\gamma}$, due to the line only, is difficult, as it has to be done in the complex visibility plane. We provide such an attempt in the velocity range where the line is clearly detected ($[-350;350]$ km/s, with line-to-continuum ratio larger than 1.05). As can be seen in the right panel of Fig. 6, the displacements are much larger (up to $\sim 1$ mas) with the largest measured at the highest velocities, and appear more spread. Nonetheless, the observed asymmetry is still consistent with the closure phase measurements that show no change through the line, within the large errors.

5. INTERPRETATION: ACCRETION OR EJECTION SIGNATURE?

As has already been discussed in previous studies, the $Br\gamma$ line could be emitted by a variety of mechanisms, such as accretion of matter onto the star, in a gaseous disk, or in outflowing matter. The spectra obtained at high spectral resolution show a double-peaked and asymmetric profile that can be interpreted in the context of the formation of optically thick lines in a dense environment with a temperature gradient.

5.1 Infalling envelope of gas?

Formation of the $Br\gamma$ line in an infalling envelope of gas can be ruled out. Considering that the line excitation temperature increases towards the star, if the line is emitted in infall of matter, or accretion flows, the profile would be double peaked but with an opposite asymmetry to what is observed (see Fig. 4, i.e. with a lower emission at redshifted velocities). In addition, in that case, the smallest extension and photocenter displacements would be expected at the highest velocities, which is in disagreement with our findings (see Figs. 4 and 6).
Figure 7. Is a jet the responsible for the photocenter displacement? Composite figure using AMBER measurements and spectro-images obtained in the [FeII] line on OSIRIS on Keck by Whelan et al. (2010). The continuum image of the ZCMa has been displayed in grey in the background of each panels. Panel (a): displacement of the photocenter measured by AMBER in different channels of the Brγ line corresponding to the right part of Fig. 6. Panels (b) to (e) are images obtained in different velocities slot by Whelan et al. (2010). The velocities channels are respectively: (b) -600 to -500 km/s, (c) -450 to -300 km/s, (d) -250 to -100 km/s and (e) -100 to 50 km/s. The blue- and redshifted emissions are located on the right side of the possible disk position angle, with the largest Brγ displacements and characteristic sizes being derived at higher velocities.

5.2 Hot layers of a gaseous disk?

The possibility that the Brγ line forms in the hot layers of the gaseous disk can also be ruled out. If one considers that the circumstellar disk surrounding the Herbig Be star is perpendicular to the large-scale jet at PA~240°, it would be expected that the velocities projected onto the line of sight cancel out along the semi-minor axis, while large spectro-astrometric displacements are seen along this axis at high velocities (Fig. 6). Apart from this, the displacements increase with velocity while Keplerian rotation should behave in the opposite way; and a phase signal is measured up to high velocities (~500 km/s), which are much larger than the expected Keplerian velocities (~100-120 km/s at ~1 AU). It therefore seems unlikely that the Brγ line is emitted in the disk.

5.3 Bipolar wind or base of a jet?

We consider that the most likely origin of the Brγ emission is a wind. Strong winds are expected to take place in massive Herbig Be and could be responsible for the Brγ emission. The double peaked and asymmetric line profile is consistent with outflowing matter emitting in optically thick lines. The visibilities and the 2D maps of the astrometric signal also support this conclusion as the blue- and redshifted emissions are located on each side of the possible disk position angle, with the largest Brγ displacements and characteristic sizes being derived at higher velocities.
Our observations suggest that the disk is slightly inclined, to allow both red- and blueshifted emissions to be seen. We may be seeing the emission from a wind partly through an optically thin inner hole in the optically thick dusty disk. Alternatively, if the inner gaseous disk were optically thick, we may be seeing the redshifted emission through a much smaller hole and via scattering on the disk surface. Fig. suggests also that the red lobe is larger than the blue lobe. Could this be due to the fact that, because of the hole in the disk, we see only a smaller part of the blue lobe on the other side of the disk? This agrees with the size of the displacements of the photocenter.

Whether the innermost disk is optically thick or not cannot be determined with our observations and no reliable estimate of the mass accretion rate exists for such high mass young stars. However, the presence of the CO overtone lines in emission (Bonnefoy et al., in prep) is indicative of a much lower mass accretion rate than those derived for FU Ors (\( \lesssim 10^{-5} M_\odot yr^{-1} \)).

At the spatial resolutions provided by the VLTI, we trace the regions close to the inner disk hole and it is therefore unsurprising that we could detect redshifted emission, while on scales of 10-100 AU, the redshifted lobe is obscured by the circumstellar disk (see Fig. Whelan et al. 2011). During this outburst, deep blueshifted absorption was detected in the Balmer lines from zero velocity to \( \sim 700 \text{ km/s} \), in addition to the absence of redshifted emission at similar velocities (Bouvier et al., in prep), supporting our conclusion that there is a strong wind in the Herbig Be.

Could our new observations be tracing the base of the jet seen at larger scale and be in fact the inner parts of the parsec scale outflow? As shown in Fig. the astrometric signal is detected at a slightly different position angle and spans a range broader than 60°, a value commonly assumed for the opening angle of the jets. At these spatial scales, it is unlikely that the jet is already collimated. The different position angle with respect to the larger scale jet (Fig. Whelan et al. 2010) could be due to jet precession. Our observations exclude a fully spherical wind since in that case no displacement would be expected between the redshifted and blueshifted emission lobes. The derived spectro-astrometric signatures favor a bipolar wind, maybe unrelated to the jet, but can not determine whether its geometry is that of a disk-wind or a stellar wind.

### 5.4 Outburst: accretion seen through massive ejection?

After detecting the same level of optical polarisation in both continuum and spectral lines along a position angle roughly perpendicular to the large-scale jet, Szeifert et al. (2010) concluded that this outburst is related to a change in the path along which the photons escape from the dust cocoon. The disappearance of the Br\( \gamma \) emission line, with respect to the continuum, after the outburst, suggests that its emission is related to the outburst. A strong mass ejection event could account for the deep blueshifted absorption features seen in the Balmer lines that are emitted close to the star as well as for the Br\( \gamma \) line emitted in outer layers of the wind. Outside the outburst, the wind disappears or is more likely to be maintained at a much smaller mass loss rate. Based on these conclusions, one can speculate about the origin of the outburst, as being driven by an event of enhanced mass accretion, similar to the EX Ors and FU Ors outbursts. In that case, this would suggest a strong link between mass accretion and ejection during the outburst, probably coupled with a magnetic field as in lower-mass young stars.

Whether the innermost disk is optically thick or not cannot be determined with our observations and no reliable estimate of the mass accretion rate exists for such high mass young stars. However, the presence of the CO overtone lines in emission (Bonnefoy et al., in prep.) suggests a much lower mass accretion rate than the ones derived for FU Ors.

### 6. CONCLUSION

We have presented spatially and spectrally resolved interferometric observations of the K-band emission in the Z CMa system. These observations were performed during the largest photometric outburst detected so far, that occurred in the innermost regions of the Herbig Be star.

We found that the Br\( \gamma \) line profile, the astrometric signal, and the characteristic sizes across the line are inconsistent with a Keplerian disk or with infall of matter. They are, instead, evidence of a bipolar wind seen through a disk hole, inside the dust sublimation radius.
The disappearance of the Br$\gamma$ emission line after the outburst suggests that the outburst is related to a period of strong mass loss. Based on these conclusions, we speculate that the origin of the outburst is an event of enhanced mass accretion, and that it does not result from a change in the system obscuration by dust. If this is valid, our results would suggest that the link between mass accretion and ejection as observed for quiescent T Tauri stars can also be at play in more massive young stars, and in high-accretion states. Is this accretion-ejection link universal, independent of the central mass?

Finally, this paper illustrates the great potential of the combination of spectro-astrometric and interferometric techniques for observing structures on micro-arcsecond scales. It may provide strong constraints on the mechanisms at play.

ACKNOWLEDGMENTS

We thank the VLTI team at Paranal, as well as R. Cesaroni, S. Antoniucci, L. Podio, P. Stee and M. van den Ancker for fruitful discussions. M.B. acknowledges funding from INAF (grant ASI-INAF I/016/07/0).

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