Birkhoff averages of Poincaré cycles for Axiom A diffeomorphisms

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Abstract

We study the time of nth return of orbits to some given (union of) rectangle(s) of a Markov partition of an Axiom A diffeomorphism. Namely, we prove the existence of a scaled generating function for these returns with respect to any Gibbs measure (associated to a Hölderian potential). As a by-product, we derive precise large deviation estimates and a central limit theorem for Birkhoff averages of Poincaré cycles. We emphasize that we look at the limiting behavior in term of number of visits (the size of the visited set is kept fixed). Our approach relies on the spectral properties of a one-parameter family of induced transfer operators on unstable leaves crossing the visited set.
1 Introduction

Let \((X, f)\) be some dynamical system preserving a probability measure \(\mu\) and pick an arbitrary Borel set \(A \subset X\) of positive \(\mu\)-measure. By Poincaré’s recurrence Theorem, \(\mu\)-almost every point \(x \in A\) comes back infinitely many times in \(A\) upon iterations of \(f\). We denote by \(r^A_n(x), n \geq 1\), the time of the \(n\)th return of \(x \in A\) to \(A\). These times are defined by induction in the following way:

\[
\begin{align*}
  r^A_1(x) &\overset{\text{def}}{=} \min\{k \geq 1, f^k(x) \in A\}, \\
  r^A_{n+1}(x) &\overset{\text{def}}{=} r^A_n(x) + r^A_1(f^{r^A_n(x)}(x)).
\end{align*}
\]

(For convenience we set \(r^A_0 \overset{\text{def}}{=} 0\).) Assuming that \(\mu\) is ergodic, the well-known Kac’s formula (virtually found in any textbook on ergodic theory) tells that \(\int r^A_1 \, d\mu = 1\). We can also introduce the Poincaré cycles of \(A\) with respect to \(x \in A\) by setting \(\tau^A_n \overset{\text{def}}{=} r^A_{n+1} - r^A_n, n \geq 0\). We obviously have \(r^A_n = \sum_{j=1}^n \tau^A_j\).

In a note [2], Birkhoff showed that

\[
\lim_{n \to +\infty} \frac{r^A_n(x)}{n} = \frac{1}{\mu(A)} \quad \text{for } \mu - \text{almost every } x \in A.
\]

(1)

This note was so overshadowed by his subsequent proof of the Ergodic Theorem for any integrable function that it escaped notice. The ‘modern’ proof of (1) can be found for instance in [16, 18] (where Birkhoff’s note is not cited) and it is a simple consequence of the Ergodic Theorem. It is natural to study the fluctuations of the convergence in (1). We can first ask for large deviations (that is of order \(O(1)\)) of this Birkhoff average of Poincaré cycles. One can also asks for log-normal fluctuations (that is of order \(O(1/\sqrt{n})\)).

As far as we know, such questions have never been investigated. It is easy to check that the random variables \(\tau^A_n\) are stationary under the conditional probability measure \(\mu_A \overset{\text{def}}{=} \mu(\cdot \cap A)/\mu(A)\). The corresponding process inherits ergodicity from ergodicity of \((X, f, \mu)\). Set in the language of dynamical systems, this means that the induced system \((A, f^\lambda, \mu_A)\) is an ergodic dynamical system [5]. The \(\tau^A_n\)’s are generally not independent. Let us mention a notable exception: When \(A\) is a state of a countable state Markov chain. A dynamical realisation of such a process is given by the so-called Gaspard-Wang map, a piece-wise linear approximation of the Manneville-Pomeau map, see e.g. [11].

The aim of the present article is to prove an accurate large deviation principle and a central limit theorem for \(r^A_n/n\) when \(f\) is an Axiom-A diffeomorphism on a Riemannian manifold \(M\), and \(\mu\) is the equilibrium state associated to a Hölder continuous potential \(\varphi\). The visited set \(A\) will be a Markov rectangle of some basic set (or a finite union of Markov rectangles).

There are a lot of recent works on return times. In most of them, “rare events” are considered, that is sets \(A_n\) such that \(\mu(A_n)\) goes to zero as \(n\) tends to infinity. Typically, \(A_n\) is a cylinder set and one looks at the rescaled return times to \(A_n\). In many dynamical systems with “sufficiently strong” mixing properties, such rescaled returns are shown to be distributed according to a
Poisson law as \( n \) goes to infinity (see [1] for a recent but not up-to-date review). For a fixed set, the moments of hitting and return times are studied in [7] in the setting of 'strongly mixing' processes. We would like to emphasize that in the present work the asymptotics are taken with respect to the number of visits to some fixed set \( A \).

Our key-result (Theorem 2.1) is the existence of a kind of “free energy” for Poincaré cycles. In probabilistic terms, we prove the existence of the scaled-generating cumulant function associated to \( r^A \). We are able to analyse the properties of this function because we show it is nothing but the logarithm of the largest eigenvalue of some one-parameter family of transfer operators. These transfer operators act on the induced system on an unstable leaf of reference crossing the set \( A \) upon consideration. This construction was done in [17] for other purposes than studying return times (namely to construct equilibrium states). Once we have this free energy for Poincaré cycles and its properties, we apply two results (that are little used) to get precise large deviation estimates and a central limit theorem for Birkhoff averages of Poincaré cycles. At a more technical level, let us notice that since return times are not continuous functions, we cannot apply the so-called contraction principle of large deviation theory [12] to the empirical measure. We can neither (directly) apply the known central limit asymptotics which are established for Lipschitz continuous functions (see [23] and references therein).

Outline of the article. In Section 2, we state our main result and its consequences. Section 3 is devoted to some preparatory notions and lemmas. The proof of the main result is given in Section 4. We first handle the case when \( A \) is a single Markov rectangle. We then show how to extend the result to a finite union of rectangles. In Section 5, we derive our large deviation and central limit theorems.

## 2 Statement of results

We refer the reader to the book of Bowen (see [3]) for the precise definitions of Axiom-A diffeomorphisms, equilibrium states and basic sets.

**Assumptions.** Throughout we assume that \( f \) is a \( C^2 \) Axiom-A diffeomorphism on a compact Riemannian manifold \( M \). Let \( \Omega \) be a basic set for \( f \) and \( \varphi : M \to \mathbb{R} \) be a Hölder continuous function. We denote by \( \mu \) the (unique) equilibrium state associated to \( \varphi \) on \( \Omega \). Finally, \( \mathcal{R} = \{ R_i \} \) denotes a finite Markov partition of \( \Omega \) (into more than one rectangle). Let \( A \subseteq \Omega \) be some finite union of atoms of the partition \( \mathcal{R} \).

The main result that we are going to prove is the following:

**Theorem 2.1.** Under the above assumptions, there exists a real number \( \alpha_0 = \alpha_0(A) > 0 \) such that for every \( \alpha < \alpha_0 \),

\[
\Psi(\alpha) \overset{\text{def}}{=} \lim_{n \to +\infty} \frac{1}{n} \log \mathbb{E}_{\mu,A} \left[ e^{\alpha r_n^A} \right] < +\infty.
\]
Moreover the map $\alpha \mapsto \Psi(\alpha)$ has the following properties:

1. The map $z \mapsto \Psi(z)$ is analytic in a complex neighborhood of $]-\infty, \alpha_0[.$
2. It is strictly convex on $]-\infty, \alpha_0[.$

We shall apply this theorem and a result due to Plachky and Steinebach [20] to get precise estimates on large fluctuations on Birkhoff averages of Poincaré cycles.

**Theorem 2.2 (Large deviations).** Under the above assumptions, we have the following estimates, for every $u \in (0, \infty)$

$$
\lim_{n \to \infty} \frac{1}{n} \log \mu_A \left\{ \frac{r^n_A}{n} \geq \frac{1}{\mu(A)} + u \right\} = \inf_{\alpha < \alpha_0} \left\{ -\left( \frac{1}{\mu(A)} + u \right) \alpha + \Psi(\alpha) \right\}
$$

and for every $0 < u < 1/\mu(A)$

$$
\lim_{n \to \infty} \frac{1}{n} \log \mu_A \left\{ \frac{r^n_A}{n} \leq \frac{1}{\mu(A)} - u \right\} = \inf_{\alpha < \alpha_0} \left\{ -\left( \frac{1}{\mu(A)} - u \right) \alpha + \Psi(\alpha) \right\}
$$

where $\alpha_0 > 0$ is the same as in Theorem 2.1.

Of course we can replace the $n$th return time, $r^n_A$, by the Birkhoff sum of Poincaré cycles, $\sum_{j=1}^n \tau^A_j$, in the previous theorem.

We also obtain a central limit theorem by using Theorem 2.1 and applying a result due to Bryc [6]. Notice that in general it is impossible to deduce a central limit theorem from a large deviation principle assuming only that the cumulant generating function is twice differentiable at the origin. Even real-analyticity is not enough (see a counterexample in [6]).

**Theorem 2.3 (Central limit theorem).** Under the above assumptions

$$
\lim_{n \to \infty} \mu_A \left\{ \frac{r^n_A - n/\mu(A)}{\sigma_A \sqrt{n}} \leq t \right\} = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{t} e^{-\xi^2/2} d\xi
$$

where

$$
\sigma_A^2 = \Psi''(0) = \lim_{n \to \infty} \frac{1}{n} \int \left( \frac{r^n_A}{\mu(A)} - \frac{n}{\mu(A)} \right)^2 d\mu_A \in ]0, +\infty[.
$$

We can replace the $n$th return time, $r^n_A$, by the Birkhoff sum of Poincaré cycles, $\sum_{j=1}^n \tau^A_j$, in the previous theorem.
Remarks

1. $\alpha_0$ has an explicit expression: it is the difference of the topological pressures of the system and the topological pressure of the system obtained by ‘removing $A$’ from the phase space, see formula (13) below.

2. We emphasize that in Theorem 2.2 we have limits, not just liminf and limsup. We could of course formulate the result in terms the Legendre transform of $\Psi$. We could also consider more general intervals of large deviations in the spirit of the Gärtner-Ellis Theorem [12].

3. It is well-known that the values assumed by $\Psi'$ when $\alpha$ ranges from $-\infty$ to $\alpha_0$ give the values of the possible deviations $u$ around the mean $1/\mu(A)$. It will be easy to check that $\Psi'(0) = \lim_{n \to +\infty} \frac{1}{n} \mathbb{E}_{\mu_A}[\tau^n_A] = \frac{1}{\mu(A)}$. Indeed,

$$
\frac{1}{n} \mathbb{E}_{\mu_A} \left( \sum_{j=1}^{\alpha} \tau^1_A \right) = \frac{1}{n} \sum_{j=1}^{\alpha} \mathbb{E}_{\mu_A}(\tau^1_A) = \mathbb{E}_{\mu_A}(\tau^1_A) = \frac{1}{\mu(A)}.
$$

The last equality is Kac formula and the fact that $\mathbb{E}_{\mu_A}(\tau^1_A) = \mathbb{E}_{\mu_A}(\tau^1_A)$ for all $j \in \mathbb{N}$ is established in [16] or [18].

4. There is another way to prove a central limit theorem close to Theorem 2.3. Let us sketch it when $A$ is a single rectangle. There is a well-known duality between the $n$th return to $A$ and the number of occurrences of $A$ up to time $n$. Let $N^n_A(x) \equiv \mathbb{I}_A(x) + \cdots + \mathbb{I}_A(f^{n-1}x)$. This is the number of visits of the orbits of $x$ to $A$ up to time $n$. A central limit theorem can be easily derived for $N^n_A$ by using [19]. Indeed, we can apply the central limit theorem given therein to the characteristic function of the one-cylinder associated to $A$ in the subshift $(\Sigma, \sigma)$ and pull it back to $(\Omega, f)$. This means that for all $t \in \mathbb{R}$

$$
\lim_{n \to \infty} \mu \left\{ \frac{N^n_A - n\mu(A)}{\sigma_A \sqrt{n}} \leq t \right\} = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{t} e^{-\xi^2/2} \, d\xi
$$

where

$$
\sigma_A^2 = \lim_{n \to \infty} \frac{1}{n} \text{Var}(N^n_A) = \lim_{n \to \infty} \frac{1}{n} \int (N^n_A - n\mu(A))^2 \, d\mu.
$$

We know that $0 \leq \sigma_A^2 < \infty$. Applying Feller’s result [13] we get [2] with

$$
\sigma_A^2 = \sigma_A^2 \mu(A)^3.
$$

(4)

Thus, proving that $\sigma_A^2 > 0$ is equivalent to prove that $\sigma_A^2 > 0$ since $A$ is a (non-trivial) Markov rectangle, that is, $0 < \mu(A) < 1$. But by [19]
we know that $\sigma_A^2 = 0$ if, and only if, $1 - \mu(A)$ is a coboundary. This cannot happen for a Markov rectangle because this would imply that for any fixed point $x \in A$, $1 = 1 - \mu(A) < 1$. Notice that by following this line of proof, we do not prove that $\sigma_A^2$ is equal to $\Psi'(0)$.

3 Preparatory lemmas

In this section, we briefly recall the relevant results from [17] which are useful for the proof of Theorem 2.1 and derive a few lemmas.

3.1 Notations

Let us recall that $\mu$ is the unique equilibrium state associated to the potential $\varphi$, that is, we have

$$h_\mu(f) + \int \varphi \, d\mu = \sup_\nu \left( h_\nu(f) + \int \varphi \, d\nu \right) = P_{\text{top}}(\varphi, \Omega)$$

where the supremum is taken over the set of $f$-invariant probability measures on $\Omega$. As usual, $h_\nu(f)$ denotes the entropy of the measure $\nu$ and $P_{\text{top}}(\varphi, \Omega)$ the topological pressure on $\Omega$ associated to the potential $\varphi$. Let $N \geq 2$ be the number of proper rectangles of the Markov partition $R$ and $A$ the $N \times N$-transition matrix defined as

$$a_{ij} = 1 \quad \text{if} \quad f^{-1}(R_j) \cap R_i \neq \emptyset$$
$$a_{ij} = 0 \quad \text{otherwise.}$$

Let $\Sigma$ be the set of sequences $x = \{x_n\}_{n \in \mathbb{Z}}$ such that for every $n$, $x_n$ belongs to $\{1, \ldots, N\}$ and $a_{x_n x_{n+1}} = 1$. If $\sigma$ denotes the shift map on $\Sigma$, there exists some canonical map $\pi$ from $\Sigma$ onto $\Omega$ such that the following diagram commutes:

$$\begin{array}{ccc}
\Sigma & \xrightarrow{\sigma} & \Sigma \\
\downarrow \pi & & \downarrow \pi \\
\Omega & \xrightarrow{f} & \Omega
\end{array}$$

As the map $\pi$ is also Hölder continuous, the map $\tilde{\varphi}$ defined by

$$\tilde{\varphi} \overset{\text{def}}{=} \varphi \circ \pi$$

is again Hölder continuous, and there exists a unique equilibrium state $\tilde{\mu}$ for the dynamical system $(\Sigma, \sigma)$ associated to the potential $\tilde{\varphi}$. Moreover $\tilde{\mu} \circ \pi^{-1} = \mu$.

The topological pressure associated to $\tilde{\varphi}$, $P_{\text{top}}(\tilde{\varphi}, \sigma)$, is equal to $P_{\text{top}}(\varphi, \Omega)$.

The cylinder set $[i_0, \ldots, i_n] \subset \Sigma$, $i_j \in \{1, \ldots, N\}$, $n \geq 0$, is the set of points $x$ such that $x_j = i_j$ (for every $0 \leq j \leq n$).

Let $g$ be the first return map in $A$: $g : A \longrightarrow A$ with $x \mapsto f^{r_A(x)}(x)$.
If $x$ is a point in $\Omega$, we denote, as usually, by $W^u(x)$, $W^u_{loc}(x)$, $W^s(x)$, $W^s_{loc}(x)$ the unstable and stable global and local manifolds. Local means that the length is equal to some expansive constant, $\varepsilon_0$. For every set $R$ of small diameter (smaller than $\varepsilon_0$) we set

$$W^i(x, R) \overset{\text{def}}{=} W^i_{loc}(x) \cap R, \text{ for } i = u, s.$$ 

We will assume that the diameter of $R$ is smaller than $\varepsilon_0$.

### 3.2 The subsystem $(F, g_F)$

For the sake of definiteness, we set $A = \pi[1]$. We denote by $F$ some fixed unstable leaf in $A$; namely we have

$$F = W^u(x_0, A)$$

for some fixed point $x_0$ in $A$. The system of local coordinates gives a projection $\pi_F$ from $A$ onto $F$. This projection is Hölder continuous. We denote by $g_F$ the map $\pi_F \circ g$. For $x$ in $\Omega$ and $x'$ in $W^s(x)$, we set

$$\omega(x, x') = \sum_{k=0}^{+\infty} \varphi \circ f^k(x) - \varphi \circ f^k(x').$$

The map $\varphi$ is Hölder continuous, and so, by contraction on the stable leaves, the previous series converges. For $x$ in $F$ we set $\omega(x) = \omega(g(x), g_F(x))$, and

$$\Phi(x) = \sum_{k=0}^{r^u_{\lambda}(x)-1} \varphi \circ f^k(x) + \omega(x).$$

This function is defined on a set of full measure with respect to any invariant measure. A simple computation gives the following lemma

**Lemma 3.1.** There exists some positive constant $C_\omega$ such that for every $x$ in $F$,

$$\|\omega\|_{\infty} \leq C_\omega.$$

The inverses branches of $g_F$ define the family of $n$-cylinders: for $x$ in $F$ we set

$$C_n(x) \overset{\text{def}}{=} f^{-r^u_{\lambda}(x)}(W^u(g^n(x), A)).$$

The $n$-cylinders are well-defined except for the points in $F$ which do not return infinitely many times in $R$ and for the points which belong to the orbit of the boundary $\partial R$ of the partition. These two sets of points will have null-measure for all the measures we are going to consider. Hence, every $n$-cylinder is a compact set and the collection of the $n$-cylinders defines a partition of $F$ (up to the boundary and points which come back less than $n$ times), which refines the partition into $(n - 1)$-cylinders. An important property is that $g^n_F(C_n(x)) = F$. 

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for every $n$-cylinder. This is a consequence of the Markov property of the partition. This allows us to define the set of preimages of some point $x$ in $F$ by $g_{\mathcal{P}}^n$, denoted by $\operatorname{Pre}_n(x)$. Hence, every $n$-cylinder contains exactly one element of $\operatorname{Pre}_n(x)$, for every $x$ in $F$. We define the Perron-Frobenius-Ruelle operator for any $x$

$$L_S(T)(x) = \sum_{y \in \operatorname{Pre}_1(x)} e^{\Phi(y) - r_S(y)ST(y)}$$

where $S$ is a real parameter. It is proved in [17] that there exists a critical value $S_c$, with $S_c \leq P_{\text{top}}(\varphi, \Omega)$, such that

$$L_S(1_{F})(x) < +\infty \text{ for every } S > S_c \text{ and for every } x \in F. \quad (6)$$

$S_c$ is of course defined as the smallest real number with this property. It is also proved that for every $x$ in $F$ and for $S = P_{\text{top}}(\varphi, \Omega)$, we still have $L_S(1_{F})(x) < +\infty$. A part of the work in the next subsection will be to prove that in fact $S_c$ is strictly smaller than the topological pressure $P_{\text{top}}(\varphi, \Omega)$ and indeed equal to the topological pressure when one removes transitions from or to $A$. Here again, the Markov property of the partition and the hyperbolic structure lead to the next lemma:

**Lemma 3.2.** There exists a positive constant $C_\varphi$ which does not depend on $S$ such that for all $x,y \in F$, $S > S_c$ and integer $n$

$$\frac{1}{C_\varphi} L^n_S(1_{F})(x) \leq L^n_S(1_{F})(y) \leq C_\varphi L^n_S(1_{F})(x).$$

There exists some quasi-metric $\eta$ on $F$ such that for every $S > S_c$, the operator $L_S$ is a quasi-compact operator on the Banach space $\mathcal{C}_\eta$ of Lipschitz-continuous function (for the quasi-metric $\eta$). We recall that the norm $\| \cdot \|_\eta$ on $\mathcal{C}_\eta$ is defined by

$$\| \phi \|_\eta = \| \phi \|_\infty + \sup_{x \neq x'} \frac{|\phi(x) - \phi(x')|}{\eta(x,x')}.$$ 

The quasi-metric $\eta$ is chosen such that the $\vartheta$-Hölder continuous functions on $F$ (where $\vartheta$ is the Hölder coefficient of $\varphi$) are $\eta$-Lipschitz continuous functions. Using the Ionescu-Tulcea & Marinescu theorem (see [13]), we get that there exists some unique probability measure $m_S$ such that

$$L_S^* (m_S) = \lambda_S m_S \quad (7)$$

where $\lambda_S > 0$. Moreover, there is a unique function $H_S$ in $\mathcal{C}_\eta$ such that

$$L_S H_S = \lambda_S H_S \quad \text{and} \quad \int H_S \, dm_S = 1.$$ 

This function satisfies

$$\frac{1}{C_0} \leq H_S \leq C_0 \quad (8)$$
where $C_0 > 0$. A consequence of Lemma 3.2 is that $C_0$ does not depend on $S$.

The measure $\nu_S$ defined by

$$d\nu_S = H_S dm_S$$

is the unique equilibrium state for $(F, g_F)$ associated to $\Phi(\cdot) - Sr_A(\cdot)$. Quasicompactness of $L_S$ means here that there exist $p = p(S)$ complex numbers of modulus one, $1 = \lambda(1), \ldots, \lambda(p)$ such that

$$L_S = \sum_{i=1}^{p} \lambda_S \lambda(i) \Psi_i + \lambda_S \tilde{\Psi}$$

(9)

where the $\Psi_i$ are linear projectors defined on $C_\eta$ with finite rank, $\tilde{\Psi}$ has a spectral radius strictly smaller than 1, and all the kernels of these operators contain the images of the others.

A crucial fact is that for $S = P_{top}(\varphi, \Omega)$, we have $\lambda_S = 1$ and $\nu_S$ is the projection on $F$ of the measure $\mu_A$.

### 3.3 Computation of $S_c$.

We can remove the set $A$ from the Markov partition to define a new subshift of finite type in the following way. For the sake of definiteness, we assume that the first line and the first column of $A$ encode the transitions to and from $A$. We denote by $A'$ the $(N-1) \times (N-1)$ matrix obtained by removing from $A$ the first line and the first column. $\Sigma'$ will denote the subset in $\Sigma$ of all sequences $x = (x_n)$ such that $a'_{x_n x_{n+1}} = 1$. For convenience we assume that the matrix $A'$ is transitive; if this is not the case, we can restrict our work to the classes of recurrence (for each class the map $\sigma$ satisfies expansiveness and specification).

The map $\tilde{\varphi}$ can be restricted to $\Sigma'$, thus the dynamical system $(\Sigma', \sigma)$ admits exactly one equilibrium state, $\tilde{\mu}'$, with topological pressure $P_{top}(\tilde{\varphi}, \Sigma')$.

**Lemma 3.3.** With the previous notations, $P_{top}(\tilde{\varphi}, \Sigma') < P_{top}(\tilde{\varphi}, \Sigma)$.

The proof can be found in [10]. Another proof based on relative entropy can be found in [8].

We can now prove the main lemma of this subsection:

**Lemma 3.4.** The critical value $S_c$, defined in (6), is equal to $P_{top}(\tilde{\varphi}, \Sigma')$.

**Proof.** To ease notations we set $P' = P_{top}(\tilde{\varphi}, \Sigma')$ throughout this proof. Let $x$ be in $F$. By Lemma 3.2 we have just to prove that for every $S > P'$, $L_S(1_F)(x) < +\infty$ and for $S = P'$, $L_S(1_F)(x) = +\infty$. We can choose $x$ such that it does not belong to the set $\bigcup_{n} f^{-n} \partial R$. Therefore $\pi^{-1}_F(x)$ is a single point $\underline{x}$ in $\Sigma$; this also holds for every $y$ in $Pre_1(x)$.

Let $n$ be some integer. The set, $Pre_1^n(x)$, of points $y$ in $Pre_1(x)$ such that $r_A(y) = n$ (which is well defined because $f^n(x)$ never belongs to the boundary of the Markov partition) is a $(\varepsilon_0, n)$-separated set of points: all these points belong to $F \subset W^{u}_{loc}(x)$ and all their images by $f^n$ belong to $W^{s}_{loc}(x)$. If we also assume
Because the matrix $A_z^\sigma\pi_n$ has only positive entries, we can assume that there exists some constant $\kappa$, independent of $S$ and $n$, such that
\[
\kappa_1 e^{(n-2)P'-nS}.
\]

Therefore the point $K$ of length $\varepsilon$ is thus proved.

We pick some $z$ in $P_{\pre}^n(x)$ and denote by $z$ its preimage by $\pi$; we also set $z = (z_n)_{n \in \mathbb{Z}}$. Hence $z_1$ and $z_{n-1}$ must be different from 1. Because of our choice of $K$, for each $W_n-2K-1$ in $E_n-2K-2$, there exist two words $w_K$ and $w_K'$ of length $K$ in $\Sigma'$ such that:

- the first letter of $w_K$ is $z_1$,
- the last letter of $w_K'$ is $z_{n-1}$,
- the word $(w_KW_{n-2K-1}w_K')$ of length $n-1$ is admissible for $\Sigma'$.

Therefore the point
\[
(\pi(W_n-2K-1)) \overset{\text{def}}{=} \ldots, z_{-1}, z_0, (w_K, W_{n-2K-1}, w_K'), z_n, z_{n+1}, \ldots
\]
(with the initial position in $z_0$) belongs to $\Sigma$, $\pi(\pi(W_n-2K-1))$ is in $P_{\pre}^n(x)$ and $\sigma^{K+1}(\pi(W_n-2K-1))$ is in the cylinder $W_n-2K-1$. The sum
\[
\sum_{y \in P_{\pre}^n(x)} e^{\Phi(y)} e^{-nS}
\]
is greater than the same sum but restricted to the $\pi(\pi(W_n-2K-1))'$s, thus by \[11\] proves that there exists some constant $\kappa_3$, which does not depend on $S$ and $n$, such that
\[
\sum_{y \in P_{\pre}^n(x)} e^{\Phi(y)} e^{-nS} \geq e^{\kappa_3 - C_\omega} e^{(n-2)P'-nS}.
\]
As $L_S(1_F)(x) = \sum_{n} \sum_{y \in P_{\pre}^n(x)} e^{\Phi(y)} e^{-nS}$, \[10\] and \[12\] prove that $L_S(1_F)(x)$ converges for every $S > P'$ and diverges to $+\infty$ for every $S \leq P'$. The lemma is thus proved.)
4 Proof of Theorem 2.1

Let us start with the case when $A$ is a single rectangle of the Markov partition. At the end of this section, we briefly sketch the modifications that are necessary to handle the case when $A$ is a finite union of rectangles.

Let us set

\[ \alpha_0 = P_{\text{top}}(\tilde{\varphi}, \Sigma) - S_c = P_{\text{top}}(\tilde{\varphi}, \Sigma) - P_{\text{top}}(\tilde{\varphi}, \Sigma'). \]  

Let $\alpha$ be in $\]0, \alpha_0[$. For convenience we will write $P$ in subscript instead of $P_{\text{top}}(\varphi, \Omega)$. We start by observing that if we pick any point on the reference unstable leaf $F$, then all points lying in its stable leaf have the same return times as this point. That is why we can reduce the problem to the unstable leaf $F$. We have

\[ I_{\text{E}} \mu_A \left[ e^{\alpha r_n A} \right] = \int_F e^{\alpha r_n A} H_P(x) dm_P(x) = \int_F \mathcal{L}_P^\alpha(e^{\alpha r_n A} H_P)(x) dm_P(x) \]

\[ = \int_F \sum_{y \in \text{Pre}_n(x)} e^{S_{\alpha}^g(y)} H_P(y) dm_P(x), \]  

(14)

We first remark that for $\alpha < \alpha_0$ we have $P - \alpha > S_c$. Hence, using lemmas 3.1-3.2, the Markov property and the properties of the function $H_P$ one obtains the existence of some constant $C$, which depends only on $\varphi$ and $f$ such that for every $x$ in $F$,

\[ e^{-C} \mathcal{L}^n_{P-\alpha}(1_F)(x) \leq \int_F \sum_{y \in \text{Pre}_n(x)} e^{S_{\alpha}^g(y)} H_P(y) dm_P(x) \leq e^{C} \mathcal{L}^n_{P-\alpha}(1_F)(x). \]  

If we integrate this double inequality with respect to the measure $m_{P-\alpha}$, we obtain the following estimate:

\[ \log \lambda_{P-\alpha} - \frac{C}{n} \leq \frac{1}{n} \log \mathbb{E}_{\mu_A} \left[ e^{\alpha r_n A} \right] \leq \log \lambda_{P-\alpha} + \frac{C}{n}. \]  

(16)

This proves that for every $\alpha < \alpha_0$

\[ \lim_{n \to +\infty} \frac{1}{n} \log \mathbb{E}_{\mu_A} \left[ e^{\alpha r_n A} \right] = \log \lambda_{P-\alpha} < +\infty. \]

Let us set $\Psi(\alpha) \overset{\text{def}}{=} \log \lambda_{P-\alpha}$. We now have to prove that the function $\Psi$ is analytic in some complex neighborhood of $\] - \infty, \alpha_0[$. Analyticity of $\Psi$ is equivalent to the analyticity of the map $S \mapsto \log \lambda_S$ in some suitable neighborhood. For that purpose, we use a theorem of perturbations due to Hennion and Hervé [14].
We first have to check that \( \mathcal{L}_S \) has \( p = p(S) \) dominating simple eigenvalues (see \([14] \), III.2). This is a consequence of ergodicity of the system \((F, g_F, \nu_S)\). In fact, the proof of the proposition 4.11 in \([4] \) can be adapted to our case; hence the projectors \( \Psi_i \) in \([4] \) have rank one, and the \( \lambda(i) \)'s all satisfy

\[
\lambda(i)^p = 1.
\]

Thus the operator \( \mathcal{L}_S \) has \( p \) dominating simple eigenvalues.

**Lemma 4.1.** Let \( \mathcal{OP}_{C_\eta} \) denote the set of linear bounded operators on \( C_\eta \). Then the map \( z \mapsto \mathcal{L}_z \) is analytic map from \( \{ z \in \mathbb{C}, \Re(z) > S_c \} \) to \( \mathcal{OP}_{C_\eta} \).

**Proof.** For any \( z \) in \( \mathbb{C} \) with \( \Re(z) > S_c \), we set (by extension):

\[
\mathcal{L}_z(\phi)(x) = \sum_{y \in \text{Pre}_1(x)} e^{\Phi(y) - r_A^1(y)} z \phi(y)
\]

which can also be written

\[
\sum_{n=1}^{+\infty} \left( \sum_{y \in \text{Pre}_1(x), \ r_A^1(y) = n} e^{\Phi(y)} \phi(y) \right) e^{-nz}.
\]

Let us introduce

\[
K_m(z)(\phi)(x) = \sum_{n=1}^m \left( \sum_{y \in \text{Pre}_1(x), \ r_A^1(y) = n} e^{\Phi(y)} \phi(y) \right) e^{-nz}.
\]

We are going to prove that the sequence \( (K_m(z)) \) converges to \( \mathcal{L}_z \), when \( m \) goes to \( +\infty \), as analytic functions from \( \Re(z) > S_c \) to \( \mathcal{OP}_{C_\eta} \). Let us fix some compact set \( \Gamma \) in \( \Re(z) > S_c \) and pick \( z \) in \( \Gamma \). Let \( \phi \) be some function in \( C_\eta \), with \( \| \phi \|_\eta = 1 \).

We want to compute

\[
\| K_m(z)(\phi) - \mathcal{L}_z(\phi) \|_\eta.
\]

There exists \( S > S_c \) such that for every \( z \) in \( \Gamma \), \( \Re(z) > S \), which proves that the series \( K_m(z)(\phi)(x) \) is normally convergent (in \( z \)) and uniformly convergent in \( x \) to \( \mathcal{L}_z(\phi)(x) \) when \( m \to \infty \). Hence \( K_m(z) \) is uniformly convergent to \( \mathcal{L}_z \) in \( \Gamma \) for the norm \( \| \cdot \|_\infty \) on \( C_\eta \). Let \( y \) and \( y' \) be two points in \( F \) with \( y' \in C_1(y) \).

Then we have

\[
e^{\Phi(y)} \phi(y) - e^{\Phi(y')} \phi(y') = e^{\Phi(y)} \phi(y) - e^{\Phi(y)} \phi(y') + e^{\Phi(y)} \phi(y') - e^{\Phi(y')} \phi(y').
\]

Therefore, the Lipschitz properties of the functions \( \phi \) and \( \Phi \) (for the quasi-symmetric \( \eta \)) and the expansion on the unstable leaves imply that there exists some constant \( C \) which depends only on \( f \) and \( \phi \) such that for every \( n \), and for every \( x \) and \( x' \) in \( F \),

\[
\sum_{y \in \text{Pre}_1(x), \ r_A^1(y) = n} \left( e^{\Phi(y)} \phi(y) - e^{\Phi(y')} \phi(y') \right) \leq C \eta(x, x'), \quad (17)
\]
where for each \( y \) in \( \text{Pre}_1(x) \), \( y' \) is the preimage of \( x' \) in the 1-cylinder \( C_1(y) \). Inequality \( 17 \) and convergence for the norm \( \| \cdot \|_\infty \) imply that the series \( K_m(z) \) is uniformly convergent to \( L_z \) in \( \Gamma \) for the norm \( \| \cdot \|_\eta \). This proves that \( z \mapsto L_z \) is analytic.

By Theorem III.8 from [14], for every \( S > S_c \), there exists some open disc in \( \mathbb{C} \) centered at \( S \), such that for every \( z \) in this disk, \( L_z \) has \( p(S) \) dominating simple eigenvalues, \( \lambda^1(z), \ldots, \lambda^p(z) \) and the maps \( z \mapsto \lambda^i(z) \) are analytic. Because the map \( \log \) is analytic in \( \mathbb{C} \setminus \mathbb{R}_- \), we can conclude that the map \( z \mapsto \log \lambda_z \) is analytic in a complex neighborhood of \( [S_c, +\infty[ \).

It remains to prove strict convexity of \( \alpha \mapsto \Psi(\alpha) \) on \( ]-\infty, a_0[ \). Convexity is obvious (by Hölder inequality). Strict convexity is equivalent to strict convexity of the map \( S \mapsto \log \lambda_S \) on \( ]S_c, +\infty[ \). We will assume that the map \( S \mapsto \log \lambda_S \) is not strictly convex and will arrive to a contradiction with the fact that \( A \) is a proper rectangle of the Markov partition.

If the map \( S \mapsto \lambda_S \) is not strictly convex, this means that its graph contains a straight line interval, say \( I \subset ]S_c, \infty[ \). In turn, this means that \( \lambda'_S/\lambda_S \) is constant on \( I \), which means that \( \lambda_S = \Lambda_A e^{\gamma_A S} \) where \( \Lambda_A, \gamma_A \) are real constants \emph{a priori} depending on \( A \). Now we invoke the unicity of analytic continuation of analytic functions to deduce that

\[
\forall z \in \mathbb{C} \text{ with } \Re(z) > S_c, \quad \lambda_z = \Lambda_A e^{-\gamma_A z}.
\]

Let us define the smallest return time in \( A \) as follows:

\[
\tau(A) = \inf\{ k \geq 1 : f^{-k}A \cap A \neq \emptyset \} = \inf\{ \tau^1_A(x) : x \in A \}.
\]

Using Lemma 3.2, we readily get for every \( S > S_c \)

\[
C_{\varphi}^{-1} \sum_{n \geq \tau(A)} \left( \sum_{y \in \text{Pre}_1(x), r^1_A(x) = n} e^{\Phi(y)} \mathbb{I}_F(y) \right) e^{-nS} \leq \Lambda_A e^{\gamma_A S} \leq C_{\varphi} \sum_{n \geq \tau(A)} \left( \sum_{y \in \text{Pre}_1(x), r^1_A(x) = n} e^{\Phi(y)} \mathbb{I}_F(y) \right) e^{-nS}.
\]

Letting \( S \to \infty \) we deduce that \( \gamma_A = -\tau(A) \).

Now, observe that

\[
\tau(A) = \Psi'(0) = 1/\mu(A)
\]

(remember Remark 3 in Section 2); we use Kač formula to obtain

\[
\tau(A) = \sum_{n = \tau(A)}^{\infty} n \mu_A\{ r^1_A = n \}.
\]
Therefore, \( r^1_A(x) = \tau(A) \) for \( \mu_A \)-almost every \( x \). Since \( \mu_A(\hat{A}) = 1 \), the topological mixing property imposes that \( \tau(A) \) must equal to one, which is absurd since the Markov partition is made of more than two rectangles.

Therefore, the function \( S \mapsto \log \lambda_S \) is strictly convex on \( [S_e, \infty[ \), so is the function \( \alpha \mapsto \Psi(\alpha) = \log \lambda_{P-\alpha} \) on \( ]-\infty, P-S_e[ = ]-\infty, a_0[ \).

**Extension to a finite union of rectangles**

We sketch how to extend the previous proof when the set \( A \) is a union \( R_{i_1} \cup \ldots \cup R_{i_k} \) of \( k \geq 2 \) rectangles (of the partition \( \mathcal{R} \)). In each rectangle \( R_{i_j} \) we pick some unstable leaf \( F_{i_j} \), and define \( F \) as the union of the \( F_{i_j} \)'s. Up to the boundary of \( \partial \mathcal{R} \), this union is a disjoint union. Thus, if we denote by \( \pi_{i_j} \) the projection from \( R_{i_j} \) onto \( F_{i_j} \) (namely \( \pi_{i_j} \), see \( \mathcal{R} \)). This defines a map \( \pi_F \) from \( A \) onto \( F \) defined as \( \bigcup F_{i_j} \) (up to some set with zero-measure for all the measures we are going to consider). Hence the map \( g_F = \pi_F \circ g \) is well defined; the Markov property of the partition \( \mathcal{R} \) allows again us to define the partitions of \( n \)-cylinders; but in this case we will not have \( g_F^P(C_n(x)) = F \), but only \( g_F^P(C_n(x)) = F_{i_j} \) for some \( j \).

The definition of the operator \( L_S \) is the same, but the value of the critical \( S, S_c \), is changing. Let \( A'' \) be the matrix obtained from \( A \) by removing the lines and the columns corresponding to the subscripts \( i_j \)'s. Let \( \Sigma'' \) be the set of sequences \( x \) such that for every \( n, a''_{n,x_{n+1}} = 1 \) and let \( P'' \) be the topological pressure for the dynamical system \( (\Sigma'', \sigma) \) associated to the potential \( \tilde{\phi} \). Then, \( S_c \leq P'' < P_{top}(\phi, \Omega) \).

Lemma 3.2 still holds but only if we pick \( x \) and \( y \) in the same \( F_{i_j} \). Accordingly, the proof of proposition 3.4 can be adapted.

Now, observe that formulas (14) and (15) are still valid, except that we have to split the integrals over the \( F_{i_j} \)'s. Therefore (16) still holds, for some positive constant \( C \).

However, it is important to notice that transitivity on \( \Omega \) implies that for every \( j \) and \( j' \), the set of points in \( F_j \) which returns infinitely many times in \( F_{j'} \) is dense. Moreover, if \( x \) is a point in \( F_{j'} \) for every \( j \), the set of preimages of \( x \) in \( F_j \) (for the map \( g_F \) is dense. Therefore exactness of \( m_S \) still holds and \( \nu_S \) is hence ergodic. We can thus again apply Theorem III.8 from [14] to get analyticity of the map \( S \mapsto \log \lambda_S \).

The proof of strict convexity is analogous to the previous case.

## 5 Proofs of Theorem 2.2 and 2.3

**Proof of Theorem 2.2** Plachky–Steinebach’s result applies once we remind that since \( \Psi \) is strictly convex and real-analytic on \( ]-\infty, a_0[ \), the function \( \alpha \mapsto \Psi'(\alpha) \) is strictly increasing on that interval.

We also observe that \( \Psi'(\alpha) \to +\infty \) when \( \alpha \to a_0 \). This is equivalent to show that \( \lambda'_S/\lambda_S \to \infty \) when \( S \to S_c \). On the other hand, \( \Psi'(\alpha) \to 0 \) when \( \alpha \to -\infty \). This can be showed by using Lemmas 3.2 and 3.4 and convexity.
Proof of Theorem 2.3. Bryc’s Theorem [6] applies. In particular, it says that the variance is equal to $\Psi''(0)$. It is easy to get formula (3) by differentiating twice $\Psi$ and using a classical factorisation. The only point to be proved is that $\sigma_A > 0$.

It is well-known that the variance can be written as follows

$$\sigma^2_A = \mathbb{E}_A((\tau^1_A)^2) - \frac{1}{\mu(A)^2} + 2 \sum_{j=2}^{\infty} \left( \mathbb{E}_A(\tau^1_A \tau^j_A) - \mathbb{E}_A(\tau^1_A)\mathbb{E}_A(\tau^j_A) \right).$$

Now we follow the proof of Proposition 4.12, p. 63, in [19] (note that Herglotz’s Theorem also applies in our setting). To this end, we just have to check that $\mathbb{E}_A(\tau^1_A \tau^j_A) - \mathbb{E}_A(\tau^1_A)\mathbb{E}_A(\tau^j_A)$ decreases exponentially fast to 0. This fact follows from the $\psi$-mixing property of the induced system (see [9]):

$$\left| \mathbb{E}_A(\tau^1_A \tau^j_A) - \mathbb{E}_A(\tau^1_A)\mathbb{E}_A(\tau^j_A) \right| \leq \sum_{p, q \in \mathbb{N}} pq \left| \mu_A(\tau^1_A = p, \tau^j_A = q) - \mu_A(\tau^1_A = p) \mu_A(\tau^j_A = q) \right| \leq C \left( \sum_{p \in \mathbb{N}} p \mu_A(B_p) \right)^2 \theta^j = C \mu(A)^{-2} \theta^j$$

where $C > 0$, $0 < \theta < 1$ and $B_p \overset{\text{def}}{=} \{ \tau^1_A = p \}$ (we used Kač formula).

We conclude that $\sigma^2_A = 0$ if and only if $\tau^1_A - 1/\mu(A)$ is a $L^2(\mu_A)$ coboundary with respect to $g$, the induced map on $A$. But if $A$ is a single Markov rectangle, this is impossible. Indeed, there exists a fixed point for $g$ which is periodic with period $\tau(A)$ for $f$ ($\tau(A)$ is defined in [15]). Reasoning as above, this leads to a contradiction with the fact that $A$ is a strict subset of $\Omega$ in measure. The case when $A$ is a (finite) union of rectangles is left to the reader.

Therefore we arrive at the conclusion that the variance $\sigma^2_A$ defined in (3) is strictly positive. Theorem 2.3 has now a complete proof.

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