INFRARED SPECTROGRAPH CHARACTERIZATION OF A DEBRIS DISK AROUND AN M-TYPE STAR IN NGC 2547

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ABSTRACT

We present 5 to 15 μm Spitzer Infrared Spectrograph (IRS) low-resolution spectral data of a candidate debris disk around an M4.5 star identified as a likely member of the ~40 Myr old cluster NGC 2547. The IRS spectrum shows a silicate emission feature, indicating the presence of warm, small, (sub)micron-sized dust grains in the disk. Of the 15 previously known candidate debris disks around M-type stars, the one we discuss in this paper is the first to have an observed mid-infrared spectrum and is also the first to have measured silicate emission. We combined the IRS data with ancillary data (optical, JHKs, and Spitzer InfraRed Array Camera and 24 μm data) to build the spectral energy distribution (SED) of the source. Monte Carlo radiation transfer modeling of the SED characterized the dust disk as being very flat (h100 = 2 AU) and extending inward within at least 0.13 AU of the central star. Our analysis shows that the disk is collisionally dominated and is likely a debris disk.

Key words: circumstellar matter – infrared: stars – open clusters and associations: individual (NGC 2547) – stars: individual (2MASS 08093547-4913033)

1. INTRODUCTION

The Spitzer Space Telescope (Werner et al. 2004) has provided the astronomical community with a wealth of information on sources with circumstellar disks, both previously known and recently discovered. The importance of having a more complete statistical sample of sources that have disks is tied in closely with our knowledge of the timescales for planet formation. For example, knowing when the inner region of the primordial or protoplanetary disk begins to clear itself of dust constrains the lower end of this timescale. Several ground-based observational near-infrared studies give evidence pointing to the clearing of the inner disk by 3–5 Myr (e.g., Haisch et al. 2001), yet the current paradigm is that planet formation occurs on timescales of 10 Myr or longer (e.g., Hillenbrand 2005). The primordial disk is essentially an accretion disk (formed as the parental molecular cloud collapsed into a star) that is still feeding the newly formed star with material. Mainly, magnetospheric accretion is responsible for the clearing of most of the gas and dust in the disk (e.g., Muzerolle et al. 2001), however, planet formation also helps with the clearing of the gaseous protoplanetary disk by the formation of planetesimals that sweep up much of the disk’s material. The formation of larger planetary bodies occurs via collisions of smaller planetesimals (lunar-mass oligarchs, see Kenyon & Bromley 2006; Chambers 2001, and references therein) and this process generates a secondary dusty disk, mostly devoid of gas, known as a debris disk, i.e., the debris disk consists mainly of dust that is produced in the collisions between the newly formed planetesimals. Since there need to be colliding planetesimals present to replenish the dust in the disk, detecting debris disk is a strong signature of an emerging planetary system.

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additional M dwarf with excess emission indicative of the presence of a debris disk. To better probe the fraction of debris disks among lower mass stars Forbrich et al. (2008) acquired an extremely deep MIPS data set (Program ID 21127; P.I.: Charles J. Lada) and discovered nine additional candidate M-dwarf debris disks. Although M-type stars are the most numerous of any type, their debris disk statistics are extremely scarce due to their faintness. Acquiring more observations of debris disks around M stars is therefore very important to improve our understanding of how these disks evolve and planet formation occurs.

In this paper we report our findings regarding the debris disk around the aforementioned member of NGC 2547, 2MASS 08093547-4913033, found by Young et al. (2004) to possess considerable excess emission at 24 μm. This source will be referred to henceforth as source 23, following the Forbrich et al. (2008) nomenclature. To determine the source’s spectral type, we acquired optical spectroscopy. We obtained Spitzer Infrared Spectrograph (IRS) observations to study this source in more detail, in particular, to determine whether the disk has silicate emission and to help constrain its spectral energy distribution (SED) modeling. The shape of silicate features gives insight into dust grain size and possibly also its crystallinity if the signal-to-noise ratio (S/N) is high enough.

The organization of the paper is as follows: Section 2 summarizes the acquisition and reduction of the optical and IRS spectroscopic data, and Section 3 presents the results obtained. We discuss the implications of our results in Section 4, and finally, we list our conclusions in Section 5.

2. OBSERVATIONS AND DATA REDUCTION

2.1. Spitzer InfraRed Spectrograph Data

The spectroscopic data of source 23 was acquired on 2006 June 19 (IRS campaign 32), through GO proposal Program ID 20750 (P.I.: Charles J. Lada). The data are comprised of low-resolution (R ~ 60–127) observations taken using the Spitzer IRS (Houck et al. 2004) Short-Low (SL) module, with a total on-source integration time of 3.2 hours. We built four Astronomical Observation Requests (AORs), each with two cycles of 240 s of ramp duration for the second order of SL (SL2: 5.2–8.7 μm) and 10 cycles of 240 s of ramp duration for the first order of SL (SL1: 7.4–14.5 μm); the AOR keys are 15096064, 15096320, 15096570, and 15096832. We used a pointing control reference sensors (PCRS) peak-up target located at (α, δ)(J2000) = (08h09m59s, 49d16′11′′) with mV = 8.14 magnitudes.

Figure 1 shows the orientation of the AORs with respect to the background in NGC 2547 during the observations.

We used the CUBE Builder for IRS Spectra Maps (CUBISM) software (Smith et al. 2007) to reduce and extract the spectrum of the source from the four AORs, for each order. The reduction began with the subtraction of the off-nod position slit from the on-source position slit for each order; this step was meant to (1) eliminate instrumental effects such as leftover background gradients from the pipeline flat-fielding, (2) subtract out the zodiacal emission, and (3) remove rogue pixels (created mostly by cosmic rays). Additional bad pixels were identified and removed. Three spectra were extracted for each order, namely, on-source spectrum, and two local background spectra on either side of the source. The SMART software (Higdon et al. 2004) was then used to merge the first and second orders, to construct an averaged local background spectrum, and also used to build the final on-source spectrum with local background subtraction.
Figure 2. Comparison of three spectra: the bottom panel shows the averaged local background spectrum (gray) and the on-source spectrum with no local background subtraction (black), while the top panel shows the on-source with local background subtraction spectrum. The PAH bands (6.25, 7.7, 8.5, 11.25, and 12.8 μm) are marked by dotted vertical lines.

Figure 3. Optical spectrum of the source 23. The inset is a zoom-in on the Hα emission line. Based on a comparison of these data to spectra of field dwarf standards, we classify this star as M4.5. The spectrum is normalized at 7500 Å.

3. RESULTS

Through a comparison to standard field dwarfs, we have measured a spectral type of M4.5 and an equivalent width $W_\lambda = 17$ Å for Hα from the optical spectrum. The spectral type we determine agrees with that obtained by Jeffries & Oliveira (2005).

The on-source, local background subtracted, IRS spectrum of Figure 2 was combined with existing multiband wavelength photometry, from the optical to MIPS 24 μm (Jeffries et al. 2004; Young et al. 2004; Forbrich et al. 2008) to build an SED, shown in Figure 4. The final spectrum shows silicate emission, however the variable background and the faintness of the source makes it very difficult to draw strong conclusions about the shape of this feature.

We model the SED using Monte Carlo radiation transfer codes developed for modeling and interpreting spectra and images of dusty accretion disks and envelopes (see Whitney et al. 2003b, 2003a, for a description of the codes). The codes include multiple anisotropic scattering and use the Bjorkman & Wood (2001) algorithm for computing radiative equilibrium dust temperatures and thermal spectra.

The presence of a silicate feature in the SED indicates that the disk contains small grains (Weingartner & Draine 2001), so we adopt the Kim et al. (1994) interstellar medium (ISM) dust model for the opacity and scattering properties (see Whitney et al. 2003b, 2003a, for a graphical representation of the wavelength dependence of this model). The grain size distribution of Kim et al. (1994) model follows a power law (slope $\sim -3$) between 0.02 μm and 0.2 μm, and then falls off exponentially.

For the object under study we adopt an axisymmetric flared disk density structure and assume the disk is heated by starlight.
Table 1
Stellar and Disk Parameters for Source No. 23

| Parameter                             | Symbol (units) | Value |
|---------------------------------------|----------------|-------|
| Adopted distance                      | $D$ (pc)       | 361   |
| Stellar temperature                   | $T_\star$ (K)  | 3089  |
| Stellar radius                        | $R_\star$ (R$_\odot$) | 0.47  |
| Stellar luminosity                    | $L_\star$ (L$_\odot$) | 0.017 |
| Stellar mass                          | $M_\star$ (M$_\odot$) | 0.11  |
| Adopted disk outer radius             | $R_{disk}$ (AU) | 300   |
| Disk mass                             | $M_{disk}$ (M$_\odot$) | $1 \times 10^{-7}$a |
| Disk inner radius                     | $R_{hole}$ (AU) | $\lesssim 0.13$ |
| Scale height at 100 AU                | $h_{100}$ (AU) | 2     |
| Disk luminosity                       | $L_{disk}$ (L$_\odot$) | 0.035 |

Note. a For a disk radius of 300 AU.

The density structure is parameterized by

$$\rho(r, z) = \rho_0 \exp \left[ -\frac{1}{2} \left( \frac{z}{h} \right)^2 \right] \left( \frac{R}{r} \right)^\alpha,$$  

(1)

where $r$ is the cylindrical radius, $z$ is the vertical scale, $\alpha$ is the radial density exponent, and $\rho_0$ is the normalization constant. The disk scale height is given by $h(r) = h_{100}(r/100 \text{ AU})^\beta$, where $h_{100}$ is the disk scale height at a radius of 100 AU and $\beta = \alpha - 1$ is the flaring power. The disk extends from an inner radius, $R_{\text{min}}$, to an outer radius $R_d$.

The lack of longer wavelength data makes it difficult to constrain disk parameters, so we fix some parameters based on previous models of irradiated disks. We adopt a surface density of the form $\Sigma(r) \sim r^{-1}$ which is typical of models of irradiated disks and fix $\beta = 1.25$ and $\alpha = 2.25$ (e.g., D’alessio et al. 1998). The SED is not very sensitive to the disk outer radius (e.g., see models in Scholz et al. 2006), so we set $R_d = 300$ AU.

For the central star we use a temperature of $T_\star = 3089$, a radius of $R_\star = 0.47$ M$_\odot$, a luminosity of $L_\star = 0.017$ L$_\odot$, and the spectrum is for a 3000 K NextGen model atmosphere (Hauschildt et al. 1999). With these star, disk, and dust parameters fixed, we constructed a small grid of SED models varying $R_{\text{min}}$, $h_{100}$ and total dust disk mass, $M_d$. The SED is somewhat degenerate in $h_{100}$ and $M_d$ and we find a good model fit using $h_{100} = 2$ AU, $M_d = 2.7(R_d/300 \text{ AU})^2M_{\text{min}}$, and disk luminosity of $L_{disk} = 0.035 L_\star$. The degeneracies in these models may be broken with longer wavelength data which is sensitive to the disk mass. We see that the silicate feature is fit rather well using the adopted ISM grain model. However, this does not rule out the possibility that grain growth has occurred in the disk and the silicate feature is produced by a population of small grains in the upper disk layers with larger grains settled toward the midplane (e.g., Scholz et al. 2007).

The best-fit model is shown in Figure 4 and the corresponding parameters are summarized in Table 1. What we can say from the modeling is that the disk is considerably lower mass and flatter than T Tauri disks. Our models also indicate that if there is an inner disk hole it is not very large, $R_{\text{min}} \lesssim 0.13$ AU, smaller than that found in other debris disks, e.g., AU Mic, typically many AU. However, the nominal model suggests that the hole is larger than the dust sublimation radius.

4. DISCUSSION

4.1. The Nature of the Disk

Although the results derived from SED modeling described above are consistent with a homologously evolved protoplanetary disk (Wood et al. 2002a), the disk appears to be devoid of gas and nonaccreting: the equivalent width of the H$\alpha$ emission line is 17 Å, which for the spectral type in question could originate in the stellar chromosphere (White & Basri 2003).

To better ascertain the nature of this disk, we first consider the relevant timescales for debris disk clearing and apply them for our source. There are four possible mechanisms for the removal of dust grains from within a debris disk (e.g., Draine & Salpeter 1979; Kruegel 2003; Grigorieva et al. 2007): (1) sublimation; (2) stellar radiation pressure; (3) stellar wind; and (4) Poynting–Roberston (P–R) drag. The sublimation temperatures for dust grains depend on their composition: amorphous carbon grains sublimate at 2000 K, whereas silicate grains sublimate at 1500 K (e.g., Preibisch et al. 1993). These temperatures are only reached at very close proximity to the star, but sublimation may be an important factor in the production of the very innermost regions of a hole, but it is not a relevant grain removal mechanism for the debris disk being discussed. The radiation pressure is also an insignificant grain removal mechanism for debris disks around M mass stars (we refer the reader to Plavchan et al. 2005, for a detailed discussion on this matter). Stellar wind pressure may be a dominant grain removal process, but only for very young and active M dwarfs with extremely high mass-loss rates ($\sim 10^{10}$ times higher than solar or $10^{-11}$ M$_\odot$ yr$^{-1}$), which is not the case for source 23. Using the mass and radius of Table 1 we calculated the escape velocity for this source to be 299 km s$^{-1}$. A Gaussian fit to the H$\alpha$ emission line profile gives a FWHM of 114 km s$^{-1}$, so it seems winds are not very energetic for this source. Regarding stellar wind drag, Plavchan et al. (2005, 2007) propose that it may be a dominant grain removal process, however, we are not able to properly evaluate this mechanism’s timescale because we do not know the exact value of the stellar mass loss rate for our source. As such, this remains an open question.

The P–R drag timescale, i.e., the time is takes for a dust grain to spiral inward into the star, is given by (Kruegel 2003)

$$\tau_{\text{P–R}} = \frac{4\pi\rho_0 c^2 r^2}{3 L_\star},$$  

(2)

where $a$ and $\rho_0$ are the radius and density of the dust grain, respectively, and $c$ is the speed of light in the vacuum.

Spitzer IRS spectra of the Deep Impact ejecta of comet 9P/Tempel 1 revealed a dust mixture composed of amorphous and crystalline silicates, amorphous carbon, water ice, and other components, with a mean density of the dust grains for this composition of 2.96 g cm$^{-3}$ (Lisse et al. 2008). A similar dust composition was found in the debris disk of HD 113766 A by Lisse et al. (2008), who observed the source using 5–35 μm Spitzer IRS data. We assume the dust in the debris disk of source 23 is similar in composition to that of HD 113766 A or comet 9P/Tempel 1 and we henceforth use the same grain density for our calculations. Equation (2) is then rewritten solely as a function of $a$ and $r$:

$$\tau_{\text{P–R}} = 0.12 \left( \frac{a}{\mu\text{m}} \right) \left( \frac{r}{\text{AU}} \right)^2 \text{Myr}.$$  

(3)

A 0.2 μm-sized dust grain at 1 AU would then take $2.4 \times 10^4$ yr to spiral onto our source. On the other hand, a 0.2 μm

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7 Using the radial dust temperature profile given by equation 28 from Kamp & Bertoldi (2000), we determine that the sublimation radius of silicate grains of average size of 0.2 μm around source 23 is 0.0045 AU or 2.0 $R_\star$.

8 The densities of amorphous carbon, graphite, olivine, forsterite, and enstatite are 1.81, 2.16, 3.32, 3.27, and 3.20 g cm$^{-3}$, respectively.
sized grain located at 41 AU would take \(\sim 40\) Myr to spiral onto the star, so small grains located beyond this radius could be leftover material from the primordial accretion disk if no other dust removal or reprocessing mechanism is at play.

Yet there is one more important physical process to be considered: grain–grain collisions. These collisions either shatter the dust grains into smaller particles (high collision velocities), or help clear the disk by triggering grain growth through agglomeration (low collision velocities); grain–grain collisions therefore re-distribute the grain size population. Grain–grain collisions will also act to reduce the vertical scale height of the disk (see e.g., Besla & Wu 2007). The best-fit SED model for our source does indeed predict a “razor-thin” or very flat disk, with a scale height of 2 AU at 100 AU (see Section 3, Table 1). The grain–grain collision timescale, \(\tau_{\text{coll}}\), is given by (Backman & Paresce 1993):

\[
\tau_{\text{coll}} \sim \frac{\tau_{\text{orb}}}{\sigma},
\]

where \(\tau_{\text{orb}}\) is the orbital period and \(\sigma\) is the dimensionless fractional surface density; the latter can be approximated by \(L_{\text{disk}}/L_\star\) (Moro-Martín et al. 2007; Moro-Martín 2007) and Equation 4 is written as

\[
\tau_{\text{coll}} \sim 0.125 \left(\frac{r}{\text{AU}}\right)^{3/2} \left(\frac{M_\odot}{M_\star}\right)^{1/2} \frac{L_\star}{L_{\text{disk}}} \text{ yr},
\]

where \(r\) is the distance of the grain from the star and \(M_\star\) is the stellar mass. Combining Equations (3) and (5), using the disk luminosity and stellar mass cited in Table 1, and a grain size of 0.2 \(\mu\)m, one finds that the collisional timescale is shorter than the P–R drag timescale for a disk radius greater than \(2.0 \times 10^{-7}\) AU. In other words, the disk is everywhere collisionally dominated. This implies that the disk is not a remnant protoplanetary disk but a bona fide debris disk since the dust grains have been reprocessed and are “second-generation.”

For disks with lower relative luminosities, the inner regions may be dominated by P–R drag, however, the outer regions are always collision dominated. For example, a disk with \(L_{\text{disk}} = 3.5 \times 10^{-3} L_\star\), with average grain sizes of 0.2 \(\mu\)m, would have an inner hole of (at least) 0.2 AU radius because in that inner region the P–R drag timescale is shorter than the grain–grain collision timescale, i.e., the dust grains would spiral in toward the star at a faster rate compared to the rate at which they are produced by grain–grain collisions.

To date, there are 15 known debris disks around M stars, namely, AU Mic (GJ 803, HD 197481), M0.5 Ve GJ 182, M1 TWA 7, M0.5 GJ 842.2, and eleven sources in NGC 2547 (Forbrich et al. 2008). AU Mic is known to have an inner hole of 17 AU and \(L_{\text{disk}}/L_\star = 6 \times 10^{-4}\) (Liu et al. 2004). Equations (3) and (5) could be used to explain an inner hole of \(\sim 0.5\) AU if the grain sizes were 0.01 \(\mu\)m, so in AU Mic’s particular system another mechanism may be clearing the dust (such as substellar or planetary companion, or stellar wind drag, Liu et al. 2004; Plavchan et al. 2005, respectively). As for GJ 182, it has 24 \(\mu\)m excess and if an inner hole exists it has a radius smaller than 1 AU according to Liu et al. (2004). TWA 7 also has 24 \(\mu\)m excess emission (Low et al. 2005) although it has not yet been possible to determine if the disk has an inner hole (Matthews et al. 2007). Lestrade et al. (2006) determined that the disk associated with GJ 842.2 has its peak emission at 140 AU although there is tentative 25 \(\mu\)m emission indicative of warm dust.

The measured disk masses of previously known debris are reported as 28 \(M_{\text{lunar}}\) (GJ 842.2 Lestrade et al. 2006), 18 \(M_{\text{lunar}}\) (TWA 7 Matthews et al. 2007), 2.11 \(M_{\text{lunar}}\) (GJ 182 Liu et al. 2004), and 0.89 \(M_{\text{lunar}}\) (AU Mic Liu et al. 2004). The fractional disk luminosities, \(L_{\text{disk}}/L_\star\), vary between \(2 \times 10^{-3}\) (TWA 7 Low et al. 2005) and \(6 \times 10^{-4}\) (Au Mic Liu et al. 2004). Regarding disk sizes, Lestrade et al. (2006) determined a radius of 300 AU for GJ 842.2 and AU Mic’s disk has a radius of 210 AU. Finally, there is a wide spread in the ages for these sources, from 200 Myr (GJ842.2; Lestrade et al. 2006) to 12 Myr (AU Mic; Strubbe & Chiang 2006). Although the disk around source 23 has an order of magnitude higher fractional luminosity than these other debris disks, it appears to be produced by a comparable amount of dust.

The disk scale height we determined from the SED modeling of the debris disk, 2 AU at 100 AU, is similar to AU Mic disk’s scale height, 2.5 AU. These scale heights are much smaller than those of T Tauri protoplanetary disks. For comparison, Wood et al. (2002b) modeled the SED of the T Tauri star HH 30 IRS, using the same Monte Carlo radiation transfer code as used in the work presented in this paper, and found that the disk scale height was 17 AU at a disk radius of 100 AU radius. Another T Tauri star whose SED was modeled using the same aforementioned code is GM Aurigae, that was found to have a scale height of 8 AU at 100 AU radius.

### 4.2. On the Silicate Feature and the Disk Mass

As mentioned above in Section 3, the silicate feature in emission indicates that there are small grains in the debris disk. Unfortunately, due to the PAH contamination and the faintness of the source, we cannot draw any conclusions regarding silicate crystallization. Although we cannot draw more quantitative conclusions from our IRS observations, to our knowledge these observations correspond to the first silicate emission detection from a debris disk around and M-type star.

To estimate the temperature of silicate grains we assume their temperature is that of a blackbody with an SED peaking at 10 \(\mu\)m, i.e., 290 K. The temperature profile obtained from the SED modeling (Section 3) indicates that the disk reaches that temperature at a radius of \(\sim 0.15\) AU (although the silicate emission could also be arising from the upper layers of the disk). The snowline, defined as the location where the temperature corresponds to that of water ice sublimation (153 K; Hayashi 1981), for this disk would be located at \(\sim 0.18\) AU. The habitable zone, \(a_{\text{HZ}}\), defined as the region in the disk where a planet can sustain liquid water at its surface, is centered roughly on a radius of 0.13 AU for our source (using the relation \(a_{\text{HZ}} = \sqrt{L_\star/L_\odot}\) from Scalo et al. 2007). To date there are seven known planet-host M dwarf stars, and these have orbits ranging from 0.02 to 2.3 AU (endl et al. 2008, and references therein). So, planet formation can indeed occur within the habitable zone or within the snowline for our source. In other words, there could be terrestrial planet formation occurring within the habitable zone, however, more data are needed to study this hypothesis in more detail.

Among the M dwarf debris disk candidates in NGC 2547, source 23 is the one that presents the biggest excess emission at 24 \(\mu\)m (Forbrich et al. 2008). Such a large excess, indicative of a relatively luminous disk, may be the result of a stochastic collision of planetary bodies within the disk (e.g., Kokubo et al. 2006) that triggered a collisional cascade (Kenyon & Bromley 2002). Such a catastrophic event happened in our own solar system, and it led to the formation of the Moon (Kleine et al. 2002).

The dust disk mass determined from the SED fitting, \(1 \times 10^{-7} M_\odot\) (or 2.7 \(M_{\text{lunar}}\)), is dependent on the distribution of the grain sizes. The fact that we detect silicate emission is an
indication that small ISM-like dust grains are indeed present in the disk. Unfortunately, no longer wavelength data are available and we are unable to further constrain the grain sizes used in our model. We can however assume that the grain sizes are smaller than those for debris disks around earlier-type stars since radiation pressure is not a significant grain removal process for M stars (e.g., the radiative blowout grain size for an A0-type is smaller than those for debris disks around earlier-type stars since we are unable to further constrain the grain sizes used in the disk. Unfortunately, no longer wavelength data are available indicating that small ISM-like dust grains are indeed present in the disk. We also showed that there is no effective grain removal mechanism for small grains in the disk, for the exception of grain growth. We believe our mass estimate is therefore a good approximation of the true disk mass for the SED fit we obtain.

5. SUMMARY

We have acquired Spitzer low-resolution spectral data from 5 to 14 μm of a M4.5 star, 2MASS 08093547-4913033 or source 23, that is a candidate member of NGC 2547, a ~ 40 Myr old cluster. Combined with previously published optical, near-infrared, IRAC, and MIPS data, we built the SED of the source and modeled several stellar and disk parameters. We briefly summarize the main results of the paper as follows.

1. We confirm excess emission longward of 8 μm for the system and report the first detection of silicate emission from a debris disk around an M-type star.
2. The silicate emission indicates that the dust grains seem to be small: micron or submicron-sized.
3. SED modeling describes the debris disk as being very flat (scale height of 2 AU at 100 AU radius), and as having a mass of ~ 2.7 lunar assuming a disk radius of 300 AU. The SED model suggests that the disk extends to within 0.13 AU of the stellar surface, outside the dust sublimation radius.
4. Comparison of the various timescales involved in removing and producing dust leads us to conclude that the disk is dominated by grain–grain collisions throughout its entire extent.
5. A stochastic collision of planetary bodies could be responsible for the higher disk luminosity of this source, and formation of terrestrial planets within the habitable zone is entirely possible for this system.
6. We confirm the spectral type of M4.5 for the source and find that its Hα emission is consistent with chromospheric activity; the Hα emission line is also narrow enough to rule out strong stellar winds.

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Facilities: Spitzer(IRIS)

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