Properties of small molecular drug loading and diffusion in a fluorinated PEG hydrogel studied by $^1$H molecular diffusion NMR and $^{19}$F spin diffusion NMR

Errol V. Mathias · Julia Aponte · Julia A. Kornfield · Yong Ba

Received: 8 June 2010 / Revised: 19 September 2010 / Accepted: 20 September 2010 / Published online: 5 October 2010
© The Author(s) 2010. This article is published with open access at Springerlink.com

Abstract Rf-PEG (fluoroalkyl double-ended poly(ethylene glycol)) hydrogel is potentially useful as a drug delivery depot due to its advanced properties of sol–gel two-phase coexistence and low surface erosion. In this study, $^1$H molecular diffusion nuclear magnetic resonance (NMR) and $^{19}$F spin diffusion NMR were used to probe the drug loading and diffusion properties of the Rf-PEG hydrogel for small anticancer drugs, 5-fluorouracil (FU) and its hydrophobic analog, 1,3-dimethyl-5-fluorouracil (DMFU). It was found that FU has a larger apparent diffusion coefficient than that of DMFU, and the diffusion of the latter was more hindered. The result of $^{19}$F spin diffusion NMR for the corresponding freeze-dried samples indicates that a larger portion of DMFU resided in the Rf core/IPDU intermediate-layer region (where IPDU refers to isophorone diurethane, as a linker to interconnect the Rf group and the PEG chain) than that of FU while the opposite is true in the PEG–water phase. To understand the experimental data, a diffusion model was proposed to include: (1) hindered diffusion of the drug molecules in the Rf core/IPDU-intermediate-layer region; (2) relatively free diffusion of the drug molecules in the PEG-water phase (or region); and (3) diffusive exchange of the probe molecules between the above two regions. This study also shows that molecular diffusion NMR combined with spin diffusion NMR is useful in studying the drug loading and diffusion properties in hydrogels for the purpose of drug delivery applications.

Keywords Fluoroalkyl double-ended poly(ethylene glycol) · Hydrogel · 5-Fluorouracil · 1,3-Dimethyl-5-fluorouracil · Drug delivery · Molecular diffusion · Spin diffusion · NMR

Introduction

Rf-PEGs (fluoroalkyl double-ended poly(ethylene glycol)) belong to a class of biocompatible and biodegradable fluorinated polymers [1–4]. Micelles formed by Rf-PEGs in water consist of a hydrophobic fluorocarbon core (Rf core) shielded by a large hydrophilic poly(ethylene glycol) shell (PEG shell). Certain combinations of the PEG chain length and fluorocarbon size lead to sol–gel two-phase coexistence and low surface erosion in water due to the high degree of cross linking of the Rf-PEG chains between the micelles [5]. These properties make the Rf-PEG potentially useful as a drug delivery depot for controlled and sustained drug delivery [5, 6]. Previously, we have used $^{19}$F $T_1$ relaxation nuclear magnetic resonance (NMR) [7] and variable temperature-dependent EPR [4] to study the Rf-PEG’s property to encapsulate hydrophobic drugs by employing a probe molecule made of free radical labeled anticancer drug, chlorambucil-tempol adduct [4, 7, 8]. Further insight into the structure of the drug-micelle assembly was obtained through molecular dynamics simulations [7]. From these studies, we also learned that the IPDU (isophorone diurethane) units which were used to link the fluoroalcohol to the PEG chain form the hydrophobic intermediate layer (IPDU intermediate layer) under...
aqueous conditions in addition to the already known Rf core and PEG shell. The IPDU intermediate layer is lipophilic in nature, and is able to hold lipophilic/hydrophobic drugs [7]. In this study, we have used molecular diffusion NMR and 19F spin diffusion NMR to study the drug loading and diffusion properties of the Rf-PEG hydrogel for a small hydrophilic anticancer drug and its hydrophobic analog.

Diffusion NMR is a noninvasive technique to study molecular diffusion in solutions [9]. Pulsed field gradient spin echo (PFGSE) NMR employs a spin echo pulse sequence together with pulsed field gradients to analyze the diffusion and other translational motion of molecules in a wide variety of media [9–12]. The diffusion coefficient and model of diffusion can be obtained. One of the first studies to focus on the diffusion of molecules in colloidal systems dates back to 1968, and was conducted by Tanner and Stejskal [13]. In their study, one artificial system of thin liquid layers (mica stack), three different kinds of plant cells (yeast, apple, and tobacco pith), and one octanol–water emulsion was studied. Useful information regarding the model of diffusion was obtained through analyses of equations for restricted diffusion. Many groups have since carried out more diffusion experiments for the studies of hydrogels [14, 15], protein binding [16] polymer chain motion in gels [17], surfactants [18, 19], oil–water emulsions [20–22], and separations [23]. To circumvent the short $T_2$ time in viscous solutions, the pulsed field gradient stimulated echo (PFGSTE) pulse sequence was developed [9, 24, 25]. Compared with the PFGSE, PFGSTE pulse sequence makes use of the possibly longer $T_1$ time to observe the NMR signal in viscous media for a longer time [26]. Another pulse sequence referred to as CONVEX (CONvection compensation/EXcitation sculpting) has been used to study the diffusion of solute molecules in PVA hydrogels [15]. This pulse sequence includes a water suppression step to avoid the interruption by the water signal during the diffusion measurement. Many other variations of NMR diffusion techniques [9, 20, 22, 27] and theoretical models have also been developed to treat hindered diffusion and restricted diffusion including diffusion in a sphere and a cylinder [22, 28, 29]. For instance, Garasanin et al. [22] examined the diffusion of the poly(dimethylsiloxane) (PDMS) fluid, its restricted diffusion inside the emulsion droplets and the Brownian diffusion of the droplets. Time-dependent diffusion NMR techniques were also developed to study the inhomogeneous diffusion of probe molecules in cross linked polymer gels [30]. The theory of solute particles undergoing restricted diffusion and the corresponding mathematical approximations useful in treating such systems have well been compiled in the review by Price [9]. A theoretical model to treat diffusion coupled with chemical exchanges was reviewed by Momot and Kuchel [31]. This model is useful for the study of drug delivery systems. A study of NMR line shapes due to diffusive exchange and relaxation process was also reported [32].

To compliment the information obtained from molecular diffusion NMR, we also performed 19F spin diffusion NMR which utilizes the dipolar interactions of abundant nuclear spins to probe spin dipolar coupled networks. Thus, molecular miscibility and domain size of spin coupled networks can be probed [33, 34].

We have used 5-fluorouracil (FU), and its hydrophobic analog, 1,3-dimethyl-5-fluorouracil (DMFU), as examples of small molecules to study the property of the Rf-PEG hydrogel as a drug delivery depot. FU belongs to a group of anticancer drugs known as antimetabolites [35, 36]. It is used for the treatment of colon cancer, rectal cancer, breast cancer, stomach cancer, and pancreatic cancer. DFMU is a hydrophobic analog of FU with both the ring amino groups being methylated [37]. FU and DFMU contain both proton and fluorine groups which make it convenient to use both the $^1$H and the $^{19}$F NMR signals to study their drug loading and diffusion characters in the Rf-PEG hydrogel.

**Experimental**

6KC6 Rf-PEG (C$_n$F$_{13}$-CH$_2$CH$_2$O-IPDU-O-(CH$_2$CH$_2$O)$_n$-IPDU-OCH$_2$CH$_2$-F$_{13}$C$_6$) was synthesized as described in previous methods [5]. The term 6KC6 refers to the PEG with MW 6,000 Daltons (6 K) and the fluoroalcohol group with six fluorinated carbon atoms (C6). IPDU refers to isophorone diurethane. In the Rf-PEG, either side of the PEG chain is linked to an Rf group through the linker IPDU. FU and DFMU with purity greater than 99% were purchased from Sigma Aldrich (USA), and their chemical structures are given in Scheme 1. All solvents used were of analytical reagent grade purchased from Fischer Scientific (USA). To prepare the gel samples, 100 mg of 6KC6 Rf-PEG and 5 mg of FU or DFMU were dissolved in 1 ml methylene chloride. The homogenous mixture was allowed to air dry completely. Then, 0.9 ml D$_2$O was added to the mixture to make the 10% 6KC6 Rf-PEG hydrogel samples.

![Scheme 1 Chemical structures of FU and DMFU](image-url)
The hydrogels were then homogenized using a Branson sonicator. The gel samples were stabilized at room temperature for more than 24 h before the NMR study. To study the drug loading properties using solid state $^{19}$F spin diffusion NMR, the FU and DMFU loaded hydrogel samples were freeze-dried with a Labconco Inc. lyophilizer.

The diffusion NMR experiments were carried out with a Bruker Avance$^{\text{TM}}$ 600 MHz system equipped with a 5-mm Diff30 water-cooled diffusion probe and a gradient amplifier capable of achieving gradient strength up to 1,200 G/cm. $^1$H PFGSTE NMR [26] and PFGSE NMR sequences [11, 12] were used. The PFGSE experiments were carried out with $\Delta=20$ ms and $\delta=1$ ms, where $\delta$ denotes the gradient duration time and $\Delta$ the distance between the two front edges of the gradient pulses. The echo intensity for each experiment was attenuated by varying the gradient strength $g$ from 0 to 1,200 G/cm. All the signal intensities relative to the corresponding $g=0$ ones were used to remove the $T_2$ effects on the NMR signal intensities. Molecular diffusions of FU and DMFU in the 6KC6 Rf-PEG hydrogel were observed through acquiring $^1$H NMR signal, and the diffusions of the Rf-PEG chains were observed through acquiring $^1$H and $^{19}$F NMR signals. All NMR experiments including those described below were carried out at 25°C.

Another kind of diffusion NMR experiment using the PFGSTET pulse sequence was also carried out in which the echo intensity was measured by varying $\Delta$ while keeping all the other parameters including $\delta$ and $g$ constant. The advantage of this technique arises from the ability to monitor the diffusion of the probe molecule for a much longer time by using a lower gradient strength. In this manner, we could see the hindrance effect of the probe molecule due to translational barriers. For short $\Delta$ values, a probe molecule may not diffuse too far to show the hindrance effect. However, when $\Delta$ is increased, this effect on the NMR signal intensity can show up by approaching nonzero echo intensity while zero echo intensity is approached for free diffusion besides the variation of the diffusion curves according to the model of diffusion. In the experiment, $\Delta$ was varied from 0 to 200 ms while the gradient strength was fixed at 75 G/cm and $\delta$ at 1 ms. To eliminate the $T_1$ and $T_2$ effects on the echoed NMR signal intensities, the signal intensity for each experiment was measured by alternating the pulse sequences with and without the gradient pulses and the ratio of the signal intensities with and without the gradients for each $\Delta$ value was then plotted against the corresponding $(\Delta-\delta/3)$ values. Samples used for this experiment include FU in D$_2$O, FU in 6K PEG D$_2$O, FU in the 6KC6 Rf-PEG hydrogel and DMFU in the 6KC6 Rf-PEG hydrogel.

In order to evaluate the locations of FU and DMFU in the Rf-PEG hydrogel, we carried out a $^{19}$F solid state spin diffusion NMR experiment for the freeze-dried FU and DMFU hydrogel samples. The drug molecules enclosed in the Rf core/IPDU intermediate-layer region may show $^{19}$F cross peaks with the peaks of the Rf group in the two-dimensional spectrum. To avoid averaging out the $^{19}$F–$^{19}$F dipolar coupling, the samples were run in a static condition.

A Doty XC4 $^1$H–$^{19}$F/X triple resonance 600 MHz solid state NMR probe was used. The two-dimensional spin diffusion pulse sequence [38] with a mixing time of 100 ms and a recycling delay of 3 s was employed.

### Results

Results for the PFGSE experiments of FU and DMFU in the Rf-PEG hydrogel are shown in Fig. 1a and b, respectively. The data are shown by plotting the natural logarithms of the normalized echo intensities $E=I_g/I_0$ (where $I_g$ is the attenuated NMR signal intensity by the pulsed field gradient and $I_0$ is the intensity without the attenuation) against $q^2 (\Delta-\delta/3)$, where $q=\gamma g \delta$. Only the experiments for $g=0–200$ G/cm were run because the NMR signal intensities decayed to lower than 10% of the original signals with gradient strengths larger than 200 G/cm. The solution for a particle’s free diffusion relating the relative echo intensity $E$ and the diffusion coefficient $D$ is given by [9]:

$$\ln(E(\Delta-\delta/3,q)) = -q^2(\Delta-\delta/3)D$$

where $\gamma$ is the nuclear gyromagnetic ratio, $D$ is the diffusion coefficient, and $q = \gamma g \delta$ (where $\delta$ and $\Delta$ have been defined in the Experimental section) [11, 12]. All the experimental points in Fig. 1 fit Eq. 1 approximately. This was, of course, done without giving a clear model of diffusion but just used as a convenient way to show the data. The calculated apparent diffusion coefficients are $6.76\pm0.09 \times 10^{-6}$ and $3.31\pm0.08 \times 10^{-6}$ (cm$^2$/s) for FU and DMFU, respectively. The apparent diffusion coefficient of FU is larger than that of DMFU in the 6KC6 Rf-PEG hydrogel. They may show the mixed effect of the diffusion coefficients for the drug molecules in the Rf core/IPDU intermediate-layer region and in the PEG–water phase (or region). The diffusive exchange of the drug molecules between the two regions may also contribute to the apparent diffusion coefficients.

To gather more experimental data in order to find the boundary effect and the exchange effect between the probe molecules in the Rf core/IPDU intermediate-layer region and in the PEG–water phase, PFGSTET diffusion experiments where $\Delta$ was varied from 0 to 200 ms while the gradient was fixed at $g=75$ G/cm and $\delta=1$ ms were carried out. The results are shown in Fig. 2, where $(\Delta-\delta/3)$ is used as the unit of the horizontal axis. Besides the samples of FU
in 6KC6 hydrogel and DMFU in 6KC6 hydrogel, experiments for FU in 6K PEG solution and FU in D_2O were also carried out for the purpose of comparison. All these sets of experimental data fit the following exponential decays well:

\[ E(\Delta - \delta/3) = y_0 + A_1 \exp[-(\Delta - \delta/3)/t_1] \]  

(2)

The corresponding parameters of \( y_0, A_1, \) and \( t_1 \) are listed in Table 1. Further discussions for the meanings of these parameters will be conducted in the Discussion section.

In order to evaluate the diffusion of the polymeric chains of the Rf-PEG itself, and thus their effects on the observed diffusions of FU and DMFU in the Rf-PEG hydrogels, we carried out \(^1\)H and \(^{19}\)F molecular diffusion NMR experiments on a sample of the Rf-PEG hydrogel (6KC6) without any probe molecules loaded in it. We found that even by using the maximum possible gradient (1,200 G/cm), it was hard to get any significant signal attenuations to estimate the diffusion coefficients for the Rf group and the PEG chains. This indicates that the hydrogel networks are quite inflexible arising from the formation of the cross linked network through the Rf-PEG chains. Therefore, we can consider the polymeric chains as being fixed in the hydrogel while only the drug molecules did the translational motions in the hydrogel.

To probe the locations of FU and DMFU in the Rf-PEG hydrogel, we carried out a two-dimensional (2D) solid state \(^{19}\)F spin diffusion NMR experiments for the freeze-dried FU and DMFU 6KC6 Rf-PEG hydrogel samples, respectively. In the 2D NMR spectrum, cross-peaks appear if two nuclei are in proximity (within some angstroms) to each other. Thus if FU or DMFU was encapsulated into the Rf core/IPDU intermediate-layer region, we may see cross peaks between the peaks of the drug fluorine groups and the Rf groups. Otherwise, no cross peak would be observed. Of course, the appearance and the intensities of the cross peaks depend on the length of the mixing time, the closeness of the atoms and the motion of the molecules. The experimental results are shown in Fig. 3, where Fig. 3a is for FU in Rf-PEG and Fig. 3b is for DMFU in Rf-PEG. The terminal CF_3 groups of the Rf core appear at −81 ppm.

![Fig. 2 Experimental NMR signal attenuations with diffusion time (\(\Delta - \delta/3\)) for DMFU in 6KC6 (squares), FU in 6KC6 (diamonds), FU in 6K PEG (upright triangle), and FU in D_2O (inverted triangles). The corresponding lines show the data fitting with the first-order exponential decays as shown in Eq. 2](image)

| Table 1 Parameters calculated using Eq. 2 for the experimental curves in Fig. 2 |
|-----------------|-----------------|-----------------|-----------------|-----------------|
|                | FU in D_2O      | FU in PEG solution | FU in Rf-PEG hydrogel | DMFU in Rf-PEG hydrogel |
| \( y_0 \)     | 0.000 ± 0.002   | 0.019 ± 0.003   | 0.078 ± 0.006   | 0.21 ± 0.01     |
| \( A_1 \)     | 1.007 ± 0.004   | 0.976 ± 0.004   | 0.92 ± 0.01     | 0.79 ± 0.01     |
| \( t_1 \) (ms\(^{-1}\)) | 40.5 ± 0.4      | 51.4 ± 0.6      | 38 ± 1          | 48 ± 2          |
and those of the CF$_2$ chains range from $-122$ to $-168$ ppm. The chemical shifts of the $^{19}$F in FU and DMFU appear at $-178$ and $-172$ ppm, respectively. The peaks in the 2D spectra are intrinsically broad due to the static condition of the samples which does not allow the dipolar interaction to be averaged out in order to detect the spin diffusions. In Fig. 3a, cross peaks only appear between the CF$_2$ and CF$_3$ $^{19}$F diagonal peaks, but no noticeable cross peak between the FU $^{19}$F peak and those of the R$_f$ groups was observed. When turned to the case of DMFU (Fig. 3b), in addition to the cross peaks between R$_f$ CF$_2$ and R$_f$ CF$_3$, the $^{19}$F nucleus of DMFU also shows cross peaks with the CF$_2$ and CF$_3$ $^{19}$F peaks. These spectra indicate that there was a significant portion of DMFU molecules in close contact with the R$_f$ core unit, while the amount of FU molecules in close contact with the R$_f$ core units might be much lower.

**Discussion**

Figure 1 showed that the apparent diffusion coefficient of FU is larger than that of DMFU in the R$_f$-PEG hydrogel. The diffusion coefficient for a dilute suspension of spherical particles is given by the Stokes–Einstein equation

$$D = R_gT/(6\pi rN_A\eta)$$  \hspace{1cm} (3)

where $R_g$ denotes the gas constant, $T$ is the temperature, $r$ is the hydrodynamic radius of a spherical particle, $N_A$ is Avogadro’s constant, and $\eta$ is the viscosity of the solution. The ratio of the apparent diffusion coefficients of FU to DMFU is 2.04. This number should not represent the ratio of the hydrodynamic radii of FU to DMFU because, as shown in Scheme 1, the molecular size difference between FU and DMFU cannot be so large even though they were not spherical and the effect of hydrations to the ratio were not considered. Thus, the different diffusion coefficients suggest that FU and DMFU diffused with different kinetic parameters and/or models in the R$_f$-PEG hydrogel. This hypothesis is also supported by the different curves shown in Fig. 2 and their corresponding parameters in Table 1. The larger $\gamma_0$ of DMFU could indicate that the diffusion of DMFU was more hindered than FU. The more hindered effect for the DMFU could have attenuated the NMR signal less. The spin diffusion results in Fig. 3 have actually shown that more DMFU resided in or was in close contact with the R$_f$ core unit than FU. In order to evaluate the hindrance effect on the drug’s diffusions in the R$_f$-PEG micelle, we have first reviewed the theoretical models of restricted diffusion [9, 10, 13, 20, 22, 29]. Starting to get some understanding into the possible model of diffusion, we first employed the simplest model used to treat diffusion in a reflecting sphere [9, 20, 22]. In this model, when the solute is in contact with a reflecting sphere or boundary, the spin is neither transported through the boundary, nor is it relaxed by the boundary. The most common approximations for this kind of diffusion include the short gradient pulse (SGP) and Gaussian phase distribution (GPD) [9, 13, 39]. Under the GPD approximation, the signal attenuation is given as an expression including the summation of the $m$th root of Bessel equation [9, 22]. In the condition of diffusion for a long period ($D\Delta \gg R^2$), this equation is reduced to the following form:

$$E = \exp\left[-q^2R^2/5\right]$$  \hspace{1cm} (4)

where $R$ denotes the maximum distance that a molecule can diffuse. The NMR signal attenuation is independent on $\Delta$.
but decreases significantly with the increase of $R$. Using Eq. 4 and the $y_0$ values in Table 1, we could obtain that the maximum distances for that FU and DMFU diffused in the Rf-PEG hydrogel were 17.8 and 13.9 µm, respectively. Of course, these results are by no means right because the boundary condition between the Rf core/IPDU intermediate-layer region and the PEG–water phase does not satisfy that defined for the restricted diffusion. From our previous study, we have learned that the diameter of the IPDU intermediate layer is about 25Å in an anhydrous condition. If the diffusion of DMFU was heavily restricted within the Rf core/IPDU intermediate layer, we would have seen a much severer restriction according to Eq. 4. Thus, besides diffusion within the Rf core/IPDU intermediate-layer region, DMFU should also have undergone longer range diffusion in the PEG–water phase. Because DMFU is much more hydrophobic than FU, more DMFU molecules could have resided in the Rf core/IPDU intermediate-layer region than the FU. (This idea has experimentally been proven by the $^{19}$F spin diffusion spectra in Fig. 3.) FU should also have behaved more or less like DMFU, but with a lesser population in the Rf core/IPDU intermediate-layer region due its higher hydrophilicity. The boundary condition between the two regions could not isolate the drug molecules in each of the regions. Thus, besides the diffusion in each of the regions, there should also be diffusive exchanges for the DMFU and the FU between the Rf core/IPDU intermediate-layer region and the PEG–water phase. In synopsis, the model of diffusion is given by the diagram of Fig. 4. Part of the FU and DMFU shown as “B” slowly diffuse (the diffusion is more hindered) in the Rf core/IPDU intermediate-layer region (seen as a smaller compartment) and part of them shown as “A” diffuse in the PEG–water phase (seen as a larger compartment) with more rapid diffusion rates. The two parts also diffusively exchange to each other. The reason for the more hindered diffusion of B is attributed to the relative immobility of the Rf chains and the IPDU units.

To theoretically describe our model of diffusion, the diffusion equation should include (1) the hindered/slow diffusion of the drug molecules in the Rf core/IPDU intermediate-layer region; (2) the relatively free/fast diffusion of the drug molecules in the PEG–water phase (although strictly to say, the PEG–water phase is not an ideally free phase); and (3) the diffusive exchange between the two regions on the boundary (or the interface) of the two regions. Bloch equations for the motions of the macroscopic nuclear magnetizations including the effect of diffusion and diffusive exchange under the time dependent pulsed field gradient can be written as [9, 31, 40, 41]:

$$\frac{\partial M_A}{\partial t} = \gamma M_A \times \vec{B}(\vec{r}, t) - \frac{M_{A,0} \gamma + M_{A,0} \gamma}{T_2} - \frac{(M_{A,0} - M_{A,0}) \gamma}{T_1} + D_N \nabla^2 M_A - k_+ M_A + k_- M_B$$

$$\frac{\partial M_B}{\partial t} = \gamma M_B \times \vec{B}(\vec{r}, t) - \frac{M_{B,0} \gamma + M_{B,0} \gamma}{T_2} - \frac{(M_{B,0} - M_{B,0}) \gamma}{T_1} + D_N \nabla^2 M_B + k_+ M_A - k_- M_B \tag{5}$$

where $M_A$ and $M_B$ denote the magnetization of the probe molecules (DMFU or FU) in the PEG–water phase and in the Rf core/IPDU intermediate-layer region, respectively, $k_+$ and $k_-$ denotes the forward and reverse exchange rate constants for A to B, and $\vec{B}(\vec{r}, t)$ denotes the magnetic field including the static magnetic field and the time dependent pulsed field gradient field. The NMR line shapes resulting from the Bloch equations where no time-dependent pulsed field gradient is included have been treated by Belton and Hills [32]. Their theoretical treatment considers how the line shape is changed due to the relaxation effect on the boundary and the diffusive exchange between the two compartments of their model system. Their treatment relies on the distinguished chemical shifts of the probe molecules in the two compartments. The distinguished chemical shifts, however, rarely exist for systems without selectively adding a paramagnetic relaxation reagent [42]. Here, we use the pulsed field gradient technique to attenuate the NMR signal intensities in order to probe the properties of diffusion for our system. The chemical shift distinction of the NMR signals of the two components is not necessary using this method. The $^{19}$F molecular diffusion experiment has shown that the Rf core was quite immobile in space. Thus, after the drug molecules entered the Rf core/IPDU intermediate-layer region, their translational freedom can be treated solely as the diffusion of the drug molecules in this region. In addition, the drug’s diffusion in this region was hindered due to the relative immobility of the Rf chains and the IPDU units. In contrast to the much smaller size of the Rf core/IPDU intermediate-layer region, as a suitable approximation, the probe molecules in the PEG–water phase can diffuse in a much longer range and the diffusion is much less hindered. It would be complicated to find the theoretical solution (more accurately to say approximation) to Eq. 5. Thus, we will discuss a few approximations specifically pertaining to the Rf-PEG hydrogel system.

@ Springer
Furthermore, all the relaxation effects in these approximations will not be included as they have been removed in the signal intensities as shown in Figs. 1 and 2.

Assuming that the diffusions of both A and B were free in each of the regions without any diffusive exchange, then Eq. 5 will be reduced to two independent Bloch equations for A and B, respectively. The solution of each of them is an exponential decay function \[ E(\Delta - \delta/3, q) = E_A(0) \exp[-q^2(\Delta - \delta/3)D_A] + E_B(0) \exp[-q^2(\Delta - \delta/3)D_B] \] (6)

This, of course, does not represent the empirical result as shown in Eq. 2, where only one exponential decay appears.

If all of the drug molecules were dissolved in the PEG-water phase, Eq. 6 will be reduced to a single exponential decay (the first term). If all of the drug molecules were restricted in the Rf core/IPDU intermediate-layer region, at a long diffusion duration \( (D\Delta \gg R^2) \), the GPD approximation would result in the solution as shown in Eq. 4. The equation to describe the diffusions of the combination of the two cases without any diffusive exchange between the two regions would be:

\[
E(\Delta - \delta/3, q) = E_A(0) \exp[-q^2(\Delta - \delta/3)D_A] + E_B(0) \exp[-q^2R^2/5] \]

where \( E_A(0) \) and \( E_B(0) \) represent the normalized signal intensities (i.e., \( E_A(0) + E_B(0) = 1 \)) corresponding to the drug’s populations in the two regions. Since \( R \) is very small for the Rf core/IPDU intermediate-layer region, the second term is approximately equal to be \( E_B(0) \). Thus, Eq. 7 becomes

\[
E(\Delta - \delta/3, q) = E_A(0) \exp[-q^2(\Delta - \delta/3)D_A] + E_B(0) \]

This could also represent the approximation, if the diffusion in the Rf core/IPDU intermediate-layer region is very slow. This equation represents the empirical solution of Eq. 2. Bear in mind that until this point, we have assumed that the diffusive exchange would not have existed. However, for diffusion limited exchange, if the diffusion coefficient \( D_B \) is much smaller than \( D_A \) and thus the exchange is very slow, this solution is a fair approximation for Eq. 5.

It has been known that for the case of free diffusions of A and B in one phase with the existence of molecular exchange between them, the solution for Eq. 5 is [31]:

\[
E(\Delta - \delta/3, q) = E_A(\Delta - \delta/3, q) + E_B(\Delta - \delta/3, q)
\begin{align*}
&= E_A(0) \frac{(k^+)^2 + (k^-)D^q^2 - k^+\Omega}{2k^-\Omega} \exp[-(\Delta - \delta/3)\frac{1}{2}(k^+ + D^q^2 + \Omega)] + E_B(0) \frac{(k^-)^2 + (k^+)D^q^2 + k^+\Omega}{2k^-\Omega} \exp[-(\Delta - \delta/3)\frac{1}{2}(k^- + D^q^2 - \Omega)]
\end{align*}
\]
where \( k^+ = k_+ + k_x \), \( k^- = k_1 - k_x \), \( q = \gamma \delta g \), \( D^+ = D_L + D_B \), \( D^- = D_L - D_B \) and \( \Omega = [(k^+ + q^2 D^+)^2 + 4k_k^2]^{1/2} \). Plot of \( E(\Delta - \delta/3, q) \) versus \( (\Delta - \delta/3) \) would show the mixed double exponential decays at a short diffusion time. With a longer diffusion time, the curve would tend to show the single exponential decay for the component with slow diffusion. This case can be related to our system by making the following conditions: (1) A is carried by the Rf core/IPDU intermediate-layer region in the PEG–water phase; (2) the diffusion of the Rf core/IPDU intermediate-layer region is zero (which has been shown by the \(^{19}F \) diffusion experiment); (3) the diffusion of B in the Rf core/IPDU intermediate-layer region is slow due to the relative immobility of the Rf chains and the IPDU units; (4) the intermediate-layer region in the PEG–water phase assuming. For the current approximation, Eq. 8 or Eq. 12 could be used to estimate the diffusion coefficients of DMFU and FU in the PEG–water phase. Using this equation, they were found to be \( 5.18 \times 10^{-6} \) and \( 6.76 \times 10^{-6} \) cm\(^2\)/s, respectively. The corresponding apparent diffusion coefficients calculated using the data of Fig. 1 are \( 3.31 \times 10^{-6} \) and \( 6.76 \times 10^{-6} \) cm\(^2\)/s, respectively. Because the \( \Delta \) value was kept small (20 ms) for the experiments done for Fig. 1, which may not allow too much time for considerable diffusive exchange, the apparent diffusion coefficients calculated from Fig. 1 should be closer to the true values than those calculated from Fig. 2 using Eq. 8 or Eq. 12.

### Conclusion

In summary, based on the results of molecular diffusion and solid state spin diffusion experiments, we have proposed the model of diffusion of small molecules in the Rf-PEG hydrogel. We have seen that a greater percentage of DMFU molecules resided in the Rf core/IPDU intermediate-layer region than that of FU. This is well correlated to the
degrees of the drug’s hydrophobicities. The drug’s diffusion occurred in the PEG–water phase through the diffusive exchange process between the Rf core/IPDU intermediate-layer region and the PEG–water phase. This provides a piece of useful information for drug delivery applications regarding the drug loading and diffusion properties of the Rf-PEG hydrogel. We would expect that small hydrophobic drugs can be held by the Rf core/IPDU intermediate-layer region to a much greater extent than hydrophilic drugs and hydrophobic drugs may tend to degrade more from the gel surface to the body together with the Rf-PEG micelles. In contrast to hydrophobic drugs, small hydrophilic drugs tend to diffuse from the PEG–water phase to the body more quickly. Our previous studies on a larger hydrophobic drug, chlorambucil-tempol adduct, shows that the drug was encapsulated in the IPDU intermediate layer [4, 7]. Besides the drug loading capacity of the IPDU intermediate layer, the Rf core may also play a role to host the DMFU molecules as can be seen from the spin diffusion spectrum in Fig. 3b. The smaller molecular size, the hydrophobicity, and the presence of the fluorine atom in the molecule may have helped to do so. We have learned that the drug loading property in the Rf-PEG hydrogel is dependent on the compatibility of the hydrophobic core-layer region of the Rf-PEG hydrogel with the size, chemical composition, and hydrophobicity of the drugs. This opens the door to design particular polymeric units to form the intermediate layer between the Rf core and the PEG shell for customized drug delivery. This study also demonstrates that the NMR techniques used in this research are useful to probe the loading and diffusion properties of drugs in polymeric micelles.

Acknowledgment This research was supported by the NSF Grant 0351848 and NSF Grant 0619147 for an NMR facility upgrade at CSULA.

Open Access This article is distributed under the terms of the Creative Commons Attribution Noncommercial License which permits any noncommercial use, distribution, and reproduction in any medium, provided the original author(s) and source are credited.

References

1. Xu B, Li L, Yekta A, Masoumi M, Kanagalingam S, Winnik MA, Zhang K, Macdonald PM, Menchen S (1997) Langmuir 13:2447–2456
2. Jeong B, Bae YH, Lee DS, Kim SW (1997) Nature 388:860–862
3. Kraft MP (2001) Adv Drug Deliv Rev 47:209–228
4. Liu X, Mao Y, Mathias EV, Ma C, Franco O, Ba Y, Kornfield JA, Wang T, Xue L, Zhou B, Yen Y (2008) J Sol-Gel Sci Technol 45:269–278
5. Giyoong T, Kornfield JA, Hubbell JA, Johannsmann D, Hogen-Esch TE (2001) Macromolecules 34:6409–6419
6. Giyoong T, Kornfield JA, Hubbell JA, Johannsmann D, Lal J (2002) Macromolecules 35:4448–4457
7. Mathias EV, Liu X, Franco O, Khan I, Ba Y, Kornfield JA (2008) Langmuir 24:692–700
8. Prabhutendolkar A, Liu X, Mathias EV, Ba Y, Kornfield JA (2006) Drug Dev 13:433–440
9. Price WS (1997) Concepts Magn Reson 9:299–336
10. Johnson CS Jr (1999) Prog Nucl Magn Reson Spectrosc 34:203–256
11. Stejskal EO, Tanner JE (1965) J Chem Phys 42:288–292
12. Stejskal EO (1965) J Chem Phys 43:3597–3603
13. Tanner JE, Stejskal EO (1968) J Chem Phys 49:1768–1777
14. Phillips RJ (2000) Biophys J 79:3350–3354
15. Regan DG, Momot KJ, Martens PJ, Kuchel PW, Poole-Warren LA (2006) Diffus Fund 4:1.1–1.18
16. Weljie AM, Yamniuk AP, Yoshino H, Isumi Y, Vogel HJ (2003) Protein Sci 12:228–236
17. Matsukawa S, Ando I (1999) Macromolecules 32:1865–1871
18. Liao Y, Basaran OA, Franses EI (2003) AIChE J 49:3229–3240
19. Söderman O, Stilbs P, Price WS (2004) Concepts Magn Reson A 23A:121–135
20. Voda MA (2009) vanDuynhoven J. Trends Food Sci Technol 20:533–543
21. Johns ML (2009) Curr Opin Colloid Interface Sci 14:178–183
22. Garasanin T, Cosgrove T, Marteaux L, Kretschmer A, Goodwin A, Zick K (2002) Langmuir 18:10298–10304
23. Viel S, Zarelli F, Caldarrelli S (2003) Proc Natl Acad Sci 100:9697–9698
24. Johnson CS Jr (1999) Prog Nucl Magn Reson Spectrosc 34:203–256
25. Tanner JE (1970) J Chem Phys 52:2523–2526
26. Cotts RM, Sun T, Marker JT, Hoch MJR (1989) J Magn Reson 83:252
27. Križ J (2004) Langmuir 20:9560–9564
28. Cornillon P, McCarthy MJ, Reid DS (1997) J Texture Stud 28:421–434
29. Söderman O, Jönsson B (1999) Macromolecules 32:86–862
30. Yamane Y, Matsui M, Kimura H, Kuroki S, Ando I (2003) Macromolecules 36:5655–5660
31. Momot KJ, Kuchel PW (2003) Concepts Magn Reson A 19A:51–64
32. Belton PS, Hills BP (1987) Mol Phys 61:999–1018
33. Klein-Douwel CH, Maas WEJR, Veeman WS, Werumeus-Buning GH, Vankan MJM (1990) Macromolecules 23:406–412
34. Yu H, Natansohn A, Singh MA, Torriani I (2001) Macromolecules 34:1258–1266
35. Blumenkranz MS, Ophir A, Claflin AJ, Hajek A (1982) Am J Ophthalmol 94:458–467
36. Joondeph BC, Peyman GA, Khooebi B, Yue BY (1988) Ophthalmic Surg 19:252
37. Kundu NG, Schmitz SA (2006) J Pharm Sci 71:935–938
38. Jenner J, Meier BH, Bachmann P, Ernst RR (1972) J Chem Phys 71:4546–4553
39. Balinov B, Jönsson B, Linse P, Söderman O (1993) J Magn Reson A 104:17–25
40. Torrey HC (1956) Phys Rev 104:563–565
41. Abragam A (1961) The principles of nuclear magnetism. Clarendon, Oxford
42. Gupta RK, Gupta P, Moore RG (1984) Annu Rev Biophys Bioeng 13:221–246