Hydrogen Permeation Through Glass

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Physical storage of gaseous hydrogen under high-pressure in glassy micro-containers such as spheres and capillaries is a promising concept for enhancing safety and the volumetric capacity of mobile hydrogen storage systems. As very low permeation through the container wall is required for storage of compressed hydrogen, development of glasses of minimal hydrogen permeability is needed. For this purpose, one has to understand better the dependence of hydrogen permeability on glass structure. The paper points out that minimizing the accessible free volume is a sound strategy to minimize hydrogen permeability. Based on previously measured and comprehensive literature data, it is shown that permeation is independently controlled by ionic porosity and network modifier content. Thus, ionic porosity in modified and fully polymerized networks can be decreased equally to the lowest hydrogen permeability among the glasses under study. Applying this concept, a drop of up to 30,000 with respect to the permeation of hydrogen molecules through silica glass is attainable.

Keywords: hydrogen storage, glass, permeability, solubility, diffusivity, ionic porosity

INTRODUCTION

Although gravimetric energy density of hydrogen gas (142 MJ kg⁻¹) is at least three times larger than that of liquid hydrocarbons like conventional fuel, it suffers from low volumetric storage capacity (Schlapbach and Züttel, 2001). Compressed hydrogen in form of onboard tanks and portable devices would offer a cost-effective and sustainable near-term solution of hydrogen storage due to its low infrastructure impact and small carbon footprint involved with the effective production of hydrogen in its gaseous state by water electrolysis or water splitting. Despite the overall attractiveness of pressurized hydrogen gas storage, there remain striking issues in selection of safe and lightweight encapsulation materials. As metals suffer from material embrittlement and cracking well-below its initial yield stress and polymers cannot offer low enough permeation rates, oxidic glasses are found to be promising candidates for cost-effective, safe, and long-term storage of pressurized hydrogen in micro containers such as hollow microspheres and capillaries (Zhevago, 2016). Hence, the usage of glass capillary arrays has received progressive attention in recent years, not least because of the well-established manufacturing process (precise redrawing tubes) but due to closed-packed (honeycomb) polymer coated structuring options in combination with a rapid refueling capability (Zhevago and Glebov, 2007; Holtappels et al., 2011; Prewitz et al., 2018).

Compressed hydrogen storage in glass micro-containers demands for low permeation through the container wall and for successful development of glass compositions for which the product of hydrogen solubility and diffusion coefficient is minimal. Besides low permeability, a strong performance of the glass during hot shaping, which is expressed mainly in the form of a sufficient workability and a small crystallization tendency, as well as during service, where chemical process robustness is important, have to be taken into account before implementing new compositions into storage applications (Prewitz et al., 2018). With respect to permeation, accessibility of free
volume of the network structure has been shown to be closely linked to the fundamental migration and solubility processes. For silica glass, predictions from noble gas solubility experiments under ambient conditions (Shackelford and Masaryk, 1978; Nakayama and Shackelford, 1990), high-pressure regimes (Shelby, 1976) and from molecular dynamics simulations (Chan and Elliott, 1991) result in a log-normal distribution of interstitial sites varying in size from 0.1 to 0.4 nm. Following the free volume approach, gas solubility and diffusion is then proportional to the fraction of the total interstitial volume that is connected to the main network and large enough for occupation (Shackelford, 2014).

Experimental studies, which are limited to fused silica and commercial borosilicate glasses (Johnson and Burt, 1922; Barrer, 1934; Altemose, 1961; Lee et al., 1962; Laska and Doremus, 1969; Laska et al., 1969; Morimoto et al., 1992; Shelby, 1996; Kurita et al., 2002), alkali-alkaline earth silicate glasses (Barton and Morain, 1970) and aluminoborosilicate glasses (Ried et al., 2014) show that isothermal permeability can vary up to three orders of magnitude. The underlying principles in the correlations between glass composition, accessible free volume and hydrogen permeation are still unclear. Only recently, a first attempt limited to hydrogen diffusivity was made by drawing a distinction between free volume reductions in modified and in fully polymerized networks (Welter et al., 2019).

The present paper picks up the threads that minimizing the accessible volume is a sound strategy to hinder migration and solubility of hydrogen molecules. Therefore, a broader data base is analyzed and a special focus is placed on the contributions in decreasing the total free volume that results from network modifying by creation of non-bridging oxygen (NBO = oxygen that is part of a tetrahedron but not shared with a second one. It forms the ionic bonds with the modifiers) and filling (“stuffing”) polymerized (NBO-free) silicate networks by introducing charge balanced network formers like aluminum and boron tetrahedra.

**Compositional Dependence**

Generally, permeation of hydrogen gas molecules through the thin wall (thickness \( d \)) of a glass capillary or hollow sphere of area \( A \) is proportional to the product of hydrogen solubility and diffusion coefficients. For the steady state (the directional flux is constant with time for \( t \geq d^2/6D \); Barrer, 1941), the \( \text{H}_2 \) permeation can be approximated using Fick’s first law (Shelby, 1996):

\[
\frac{dQ}{dt} = \frac{A \Delta p}{d} P = \frac{A \Delta p}{d} SD, \tag{1}
\]

where \( dQ/dt \) is the permeation rate (mol s\(^{-1}\)), \( \Delta p \) is the pressure difference (Pa) between inner and outer surface of the wall and \( S \), \( D \), \( P \) are the coefficients of hydrogen solubility (mol m\(^{-3}\) Pa\(^{-1}\)), diffusion (m\(^2\) s\(^{-1}\)) and permeation (mol m\(^{-1}\) Pa\(^{-1}\) s\(^{-1}\)), respectively. Equation (1) assumes Henry’s law (physical dissolution), and that coefficients are independent on concentration. Solubility, diffusivity and permeability depend strongly on temperature. Within narrow temperature ranges, permeability obeys an Arrhenius type dependence, which can be expressed as:

\[
P = P_0 \exp \left( - \frac{E_p}{RT} \right), \tag{2}
\]

where \( P_0 \) is the maximum hydrogen permeation at infinite temperature, also named permeation constant (mol m\(^{-1}\) Pa\(^{-1}\) s\(^{-1}\)), \( E_p \) is the activation energy (J mol\(^{-1}\)) for permeation, \( T \) is the temperature (K) and \( R \) is the gas constant (8.314 J mol\(^{-1}\) K\(^{-1}\)). If broader temperature ranges are analyzed, a slight temperature dependence of the activation energy of Equation (2) becomes evident. Under these conditions, temperature dependence of gas permeability in glass is best described by Shelby and Keeton (1974) and (Shelby, 1977a):

\[
P = P_0 \exp \left( - \frac{E_p}{RT} \right) \tag{3}
\]

Tsugawa et al. (1976) introduced a first empirical correlation between hydrogen permeation and glass composition. They were following a similar compositional approach for helium permeation (Altemose, 1961), which was based on the finding that for commercial borosilicate glasses \( E_p \) decreases with increasing fraction of network former, whereas \( P_0 \) is constant with composition. Under these constraints, Equation (2) is given by (Tsugawa et al., 1976):

\[
P = 8.1 \times 10^{-14} \exp \left( - \frac{17,330 - 127.8G}{T} \right), \tag{4}
\]

where \( P \) is in units of mol m\(^{-1}\) Pa\(^{-1}\) s\(^{-1}\) and \( G \) is the total molar fraction (%) of network forming SiO\(_2\), P\(_2\)O\(_5\), B\(_2\)O\(_3\), and Al\(_2\)O\(_3\). Analysis of the available data shows that Equation (4) can only serve as an approximate of the hydrogen permeation through silicate glasses (Figure 1). Worthy of note here is that \( P_0 \) and \( E_p \) differ considerably among compositions of similar network former fraction. This finding gives clear evidence that permeation, which reflects both, solubility and diffusivity, is not solely governed by \( G \). Although \( P_0 \) strongly varies, Figure 1 indicates that \( P_0 \) and \( E_p \) tend to decrease with increasing \( G \).

In order to assess to which extent solubility \( S \) and diffusivity \( D \) could interfere in different types of glasses as \( P = SD \) and \( E_p = E_D + E_S \), the activation energies for diffusivity and solubility, \( E_D \) and \( E_S \), are plotted in Figure 2. For that purpose, data of hydrogen, deuterium, and noble gases were analyzed for glasses containing non-bridging oxygen (NBO) and those, which are nominal NBO-free. Therefore, the degree of depolymerization in glasses was determined for each composition by the ratio of non-bridging oxygen over tetrahedra cations (NBO/T). NBO/T is calculated considering the molar content of network modifier oxides, e.g., (Na\(_2\)O + CaO + K\(_2\)O + MgO) and all network forming oxides, e.g., (SiO\(_2\), Al\(_2\)O\(_3\), B\(_2\)O\(_3\)). For commercial aluminoborosilicate glasses, we can assume that all aluminum and boron are four-fold coordinated and stabilized by modifier ions. Then, NBO/T is given as:

\[
\frac{\text{NBO}}{T} = 2 \frac{\left( \text{Na}_2\text{O} + \text{K}_2\text{O} + \text{MgO} + \text{CaO} \right) - \left( \text{Al}_2\text{O}_3 + \text{B}_2\text{O}_3 \right)}{\left( \text{SiO}_2 + 2\text{Al}_2\text{O}_3 + 2\text{B}_2\text{O}_3 \right)}, \tag{5}
\]
As a result, Figure 2 reveals different behavior of NBO-bearing and NBO-free glasses. In case of NBO-bearing glasses, the activation energy of solubility, $E_S$, ranges between $-5$ and $+20$ KJ mol$^{-1}$ and that of diffusion, $E_D$, between $+25$ and $+50$ KJ mol$^{-1}$. No correlation between $E_D$ and $E_S$, however, is seen. One can state that more activation energy is needed to solve gas atoms or molecules as in case of NBO-free glasses. For the latter, $E_S$ is mostly negative varying between $-5$ and $-20$ KJ mol$^{-1}$, whereas $E_D$ more strongly varies between $+20$ and $+120$ KJ mol$^{-1}$. In contrast to NBO-bearing glasses, for (NBO/T = 0) $E_D$ and $E_S$ correlates according to

$$E_S = 2.773 - 0.197 \times E_D$$

($E_S$ and $E_D$ in kJ mol$^{-1}$), indicating that similar fundamental (thermally activated) processes for occupation and migration of gases are active. However, a negative $E_S$ means that $S$ decreases with temperature and that hydrogen molecules are likely to be dissolved in NBO-free glasses. This would imply stronger glass-gas interactions. The decrease in $E_S$ with increasing $G$ can also explain the similar trend shown in Figure 1B.

**DEPENDENCE ON FREE VOLUME**

Volumetric properties of glasses can be accessed from density and chemical composition. Firstly, one can calculate the molar volume $V_M$ (based on 1 mol oxygen atoms) as

$$V_M = \frac{\sum x_i M_i}{n} \times \frac{1}{\rho},$$

where $x_i$, $M_i$, and $n_i$ are being, respectively, the mol fraction, the molar mass of the oxide $i$ and the number of oxygen atoms in the
oxide \(i\) with \(A_iO_k\) chemical formula. Secondly, \(V_M\) can be utilized to determine the total free volume \(V_{\text{free}}\) and ionic porosity \(IP\) (Carroll and Stolper, 1993). \(V_{\text{free}}\) and \(IP\) are, respectively, the part and fraction of the molar volume not occupied by individual atoms, which are given by:

\[
V_{\text{free}} = V_M - \sum_i x_i V_i \tag{8}
\]

and

\[
IP = 1 - \frac{1}{V_M} \sum_i x_i V_i \tag{9}
\]

where \(V_i = 4\pi N_A (r_A^3 + k r_O^3)/3\) is the ionic volume the oxide \(i\) with \(A_iO_k\) chemical formula. \(N_A\) is the Avogadro constant and \(r_A, r_O\) are ionic radii using, 66 pm (Li\(^{+}\)), 102 pm (Na\(^{+}\)), 151 pm (K\(^{+}\)), 166 (Rb\(^{+}\)), 188 pm (Cs\(^{+}\)), 118 pm (Sr\(^{+}\)), 142 pm (Ba\(^{+}\)), 74 pm (Zn\(^{+}\)), 140 pm (Pb\(^{+}\)), 11 pm (B\(^{+}\)), 39 pm (Al\(^{+}\)), 26 pm (Si\(^{+}\)), 17 pm (P\(^{+}\)), and 135 pm (O\(^{+}\)) (Shannon, 1976).

**Figure 3A** shows the isothermal permeability of hydrogen/deuterium molecules at 473 K as a function of the ionic porosity at that temperature. In order to allow for analysis on a broader data basis, permeation of helium atoms was included (**Figure 3B**). Overall, it is very clear that permeation of helium and hydrogen seems to follow the expected trend, i.e., a positive correlation with ionic porosity is evident. On the other hand, a network of high connectivity seems to be more effective in reducing permeability as the data points of fully polymerized network structures (NBO-free) are mainly located on the right hand side of the hatched areas of **Figure 3**.

Further, one can analyse that permeability is higher for helium atoms than for hydrogen molecules. Both data sets are shifted in by a factor of approximately 36, with regard to the mean value
for silica glass (labeled by a red star in Figure 3). It is quite understandable that this difference can be traced back to the difference in the kinetic diameter of hydrogen, which is about 11% larger than that of helium $d_{\text{H}_2} = 298$ pm (Chambers et al., 1998), $d_{\text{He}} = 260$ pm (Matteucci et al., 2006).

Furthermore, the data sets of Figure 3 seem to follow self-similar pattern (hatched areas), which will be utilized to access volumetric dependencies with the following master-curve concept. For that reason, dimensionless ratios (the so-called "reduced permeability") were calculated for isothermal conditions by dividing permeability through the respective value $[2.2 \times 10^{-17}$ mol m$^{-1}$ Pa$^{-1}$ s$^{-1}$ ($\text{H}_2$/$\text{D}_2$) and $8.0 \times 10^{-16}$ mol m$^{-1}$ Pa$^{-1}$ s$^{-1}$ (He)] for silica glass. Figure 4 shows that reduced permeability declines in fully polymerized (NBO-free) glasses and depolymerized (NBO-bearing) glasses by up to four and a half orders of magnitude. It becomes clear that stuffing the network by charge compensating cations in fully polymerized networks is a viable option for considerably reducing hydrogen permeability. Particularly, these glasses are favorable when comparing their process properties as they offer a smaller crystallization tendency and a higher workability than NBO-bearing compositions.

Figure 5 shows the dependence of reduced permeability on ionic porosity for fully polymerized glasses and depolymerized glasses in separate figures. Grouping the NBO-free glasses by network former (Figure 5A) reveals that charge balancing of tetrahedral boron is less effective in decreasing reduced permeability. In particular, technical borosilicate glasses are only capable to decrease the permeability of silica by a factor < 100. In contrast, when tetrahedral aluminum is stabilized by modifier ions, available data show that reduced permeability is strongly declined. Factors up to $10^{4.5}$ in reduction are found for lithium and sodium aluminosilicate glasses. It is interesting to note that stabilization of mixed network formers (aluminoborosilicate glass) lead to an intermediate position between borosilicate and aluminosilicate compositions. On the other hand, if network modifiers are present in excess, they depolymerize the network and compensate charges of NBOs instead of tetrahedral coordinated aluminum and boron. In this case, reduction of free volume and with it the drop in reduced permeability depends on the nature of the modifier cation (Figure 5B). High-field strength cations such as Li$^+$, which strongly attract NBOs, lead to a relative low reduction in ionic porosity but to a relative strong decrease in permeation, whereas low-field strength cations such as K$^+$, Rb$^+$ and Cs$^+$, which are less effective to attract NBOs, result in stronger reduction of free volume but less decrease in permeation. The compositional dependence becomes clearer if one follows the iso-NBO/T lines. For NBO/T = 1 (in technology called disilicate compositions) on can find that lithium ions reduce permeation by a factor of 30,000, sodium ions by 1,800, and potassium ions only by 250 with respect to the permeation through silica glass. Figure 5B shows further that for Na$^+$-modified glasses, which includes the mixed modifier float glass of high technical importance, the dependence of reduced permeability on ionic porosity is only slightly weaker than for binary lithium silicates glasses.
CONCLUSIONS

The analysis of available data show that assuming a simplistic effect of either ionic porosity or network modifier content fails to predict hydrogen permeability in silicate glasses. Fully polymerized (NBO-free) glasses, which are stuffed with $\text{Li}^+$ and $\text{Na}^+$ ions for charge balancing network forming tetrahedral aluminum as well as depolymerized (NBO-bearing) lithium and sodium silicate glasses, both ranging between 50 and 51.5% ionic porosity, seem to offer low hydrogen permeability that is up to 30,000 times lower than that of silica glass. This finding indicates that, the local network connectivity may control hydrogen permeability more strongly than ionic porosity.

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DATA AVAILABILITY STATEMENT

The datasets generated for this study are available on request to the corresponding author.

AUTHOR CONTRIBUTIONS

TW, UM, and SR produced the manuscript. JD and RM contributed to the critical revision and direction of the manuscript.

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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