The Carina Project. IV. Radial Velocity Distribution

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ABSTRACT. We present new and accurate radial velocity (RV) measurements of luminous stars of all ages (old horizontal branch, intermediate-age red clump, and young blue plume, as well as red giants of a range of ages: 20.6 ≤ V ≤ 22) in the Carina dwarf spheroidal galaxy, based on low-resolution spectra collected with the FORS2 multiobject slit spectrograph at the VLT. This data set was complemented by RV measurements based on medium- and high-resolution archive spectra of brighter (V ≲ 20.6) Carina targets collected with the GIRAFFE multiobject fiber spectrograph at the VLT. The combined sample includes more than 21,340 individual spectra of ≈2000 stars covering the entire body of the galaxy. The mean (RV) = 220.4 ± 0.1 km s\(^{-1}\)) and the dispersion \(σ ≈ 11.7 \pm 0.1 \text{ km s}\(^{-1}\)) of the RV distribution of candidate Carina stars (∼1210 objects, 180 ≤ RV ≤ 260 km s\(^{-1}\), 4σ) agree quite well with similar measurements available in the literature. To further improve the statistics, the accurate RV measurements recently provided by Walker et al. were also added to the current data set. We ended up with a sample of ∼1370 RV measurements of candidate Carina stars that is ∼75% larger than any previous Carina RV sample. We found that the hypothesis that the Carina RV distribution is Gaussian can be discarded at 99% confidence level. The mean RV across the body of the galaxy varies from ∼220 km s\(^{-1}\) at a distance of 7′ (∼200 pc) from the center to ∼225 km s\(^{-1}\) at 13′ (∼400 pc, 6σ level) and flattens out to a constant value of ∼221 km s\(^{-1}\) at larger distances (600 pc, 4σ level). Moreover, and even more importantly, we found that in the Carina regions where the mean RV is smaller, the dispersion is also smaller, and the RV distribution is more centrally peaked (i.e., the kurtosis attains larger values). The difference in mean RV is more than 4 km s\(^{-1}\) (9σ level) when moving from east to west and more than 3 km s\(^{-1}\) (7σ level) when moving from north to south. The RV gradient appears to be in the direction of the Carina proper motion. However, this parameter is affected by large uncertainties to reach a firm conclusion. There is evidence of a secondary maximum in RV across the Carina center (|D| ≤ 200 pc). The occurrence of a secondary feature across the Carina center is also supported by the flat-topped radial distribution based on the photometric catalog. These findings are reminiscent of a substructure with transition properties that are already detected in dwarf ellipticals, and they call for confirmation by independent investigations.

Online material: color figures

1 Based on spectra collected with the spectrograph FORS2 available at the ESO Very Large Telescope, Cerro Paranal, (072.D-0671(A) PI: Bono; 078.B-0567(A) PI: Thévenin).
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1. INTRODUCTION

Dwarf galaxies are fundamental laboratories to investigate the influence of the environment on star formation and on chemical evolution in stellar systems that are several orders of magnitude smaller than giant galaxies. Empirical evidence indicates that in the Local Group (LG), the dwarf ellipticals (dE), and the dwarf spirals (ds, M33-like) we are facing three different dwarf morphological types. The dwarf spheroidals (dSphs) show either single (Cetus, Monelli et al. 2010a; Tucana, Monelli et al. 2010b) or multiple star-formation events (Carina, Bono et al. 2010), are spheroidal in morphology, and are devoid of neutral hydrogen (Bouchard et al. 2005). The dwarf irregulars (dIs) host a mix of old and young stellar populations, but they have recently experienced an intense SF episode (Sanna et al. 2009; Cole 2009); moreover, they have irregular morphology and host a significant fraction of neutral hydrogen. The transition dwarfs have properties between the dSphs and the dIs. The observational scenario concerning the dwarfs in the LG was recently enriched by the discovery of ultrafaint dSphs in the SDSS (Zucker et al. 2006; Tolstoy et al. 2009; Wyse 2010), but their properties still need to be investigated in detail. Although the morphological classification appears robust, we still lack firm quantitative constraints concerning the evolution/transition between different dwarf morphological types. Kormendy (1985, 1987) found a well-defined dichotomy between ellipticals (Es) and spheroidals (Sphs) (Kormendy et al. 2009; Graham & Guzmán 2003). Early-type galaxies are distributed along a sequence moving from cD to dEs, while spheroidals are distributed along a sequence that overlaps with spirals and dwarf irregulars (see Fig. 1 in Kormendy et al. 2009). This empirical dichotomy seems to suggest that E and Sphs galaxies are stellar systems that underwent different formation and evolution processes. In particular, the Sphs might be either spirals or irregulars that lost their gas or transformed it into stars. However, there is no general consensus concerning the E-Sph dichotomy, since the correlation between the shape of the brightness profile and the galaxy luminosity is continuous when moving from Es to Sphs (Jerjen & Binggeli 1997; Gavazzi et al. 2005; Ferrarese et al. 2006).

Theoretical and empirical uncertainties affect our understanding of several empirical correlations. More than 20 years ago it was suggested by Skillman et al. (1989) that the luminosity of dwarfs is correlated with their mean metallicity. However, it is not clear whether this correlation is linear over the entire metallicity range or shows a plateau in the metal-poor regime (Helmi et al. 2006; Mateo et al. 2008; Kirby et al. 2008). More recently, it was suggested by Woo et al. (2008) that dEs/dSphs in the LG at fixed total visual luminosity show larger mass-to-light (M/L) ratios (see their Fig. 1). At a fixed total stellar mass, the same systems are more metal-rich than dIs (see their Fig. 10 and Mateo 1998).

Precise and homogeneous photometric and spectroscopic data are required to address the preceding open issues concerning the formation and evolution of these systems. The advent of multiobject spectrographs for the 8–10 m class telescopes provided the unique opportunity to collect sizable samples of low-, medium-, and high-resolution spectra covering the entire body of nearby dwarf galaxies. By using hundreds of radial velocity (RV) measurements, it has been suggested by Kleya et al. (2004) that Sextans dSphs host kinematic substructures; in particular, they found that the RV dispersion of the stellar population located across the center of this system is close to zero and increases outside the core. A different kinematic status has also been suggested by Battaglia et al. (2008) for the two distinct stellar populations in the Sculptor dSph.

Moreover, evidence of a falloff in the velocity dispersion at large galactic distances was suggested in Sextans, Draco, and Ursa Minor (Kleya et al. 2003; Wilkinson et al. 2004, 2006). However, using large and homogeneous samples of RV measurements based on high-resolution spectra collected with the Michigan/MIKE fiber system (MMFS) at Magellan (Walker et al. 2007, hereafter W07), Walker et al. (2009) found that the RV profile in several dSphs (Carina, Draco, Leol, LeolII, and Ursa Minor) are flat at large projected radii (R > 1 kpc), while Fornax, Sculptor, and Sextans show a gentle decline.

The observational scenario is far from being settled; indeed, using the same RV measurements provided by W07, but a different algorithm to reject tidally stripped stars, Łokas et al. (2008 and Łokas 2009) found a significant decline in the RV dispersion profiles of Carina, Fornax, Leol, Sculptor, and Sextans. Furthermore, it was also suggested that kinematic status of Carina (Muñoz et al. 2006) and Bootes (Belokurov et al. 2006) might be disturbed by the Milky Way. In this context, Carina plays a key role because it is relatively close (μ = 20.10 mag), shows at least two well-separated SF episodes (t = 1–6 and 12 Gyr), and it is metal-poor (<[Fe/H] > ~ − 1.7 dex).

Our group is involved in a long-term project on the evolutionary properties, variable stars, and kinematic and abundances of stellar populations in the Carina dSph. This is the fourth article of a series, and it is focused on the radial velocity distribution of candidate Carina stars. The structure of the article is as follows. In § 2, we discuss in detail the different spectroscopic data sets we collected for this experiment, together with the approach we adopted to reduce the data. Section 3 deals with radial velocity measurements and with the observational strategy we adopted to validate radial velocities based on low-, medium-, and high-resolution spectra. In § 4 we lay out the selection criteria adopted to identify candidate Carina stars and their radial velocity distribution. In § 5, we investigate the velocity distribution, the velocity dispersion, and the kurtosis of candidate Carina stars as a function of the projected radial distance. In that section, we also address the similarity between spectroscopic and photometric radial distributions. Finally, in § 6 we summarize the results of this investigation and briefly outline the formation and evolution of these systems. The advent of multiobject spectrographs for the 8–10 m class telescopes provided the unique opportunity to collect sizable samples of low-, medium-, and high-resolution spectra covering the entire body of nearby dwarf galaxies. By using hundreds of radial velocity (RV) measurements, it has been suggested by Kleya et al. (2004) that Sextans dSphs host kinematic substructures; in particular, they found that the RV dispersion of the stellar population located across the center of this system is close to zero and increases outside the core. A different kinematic status has also been suggested by Battaglia et al. (2008) for the two distinct stellar populations in the Sculptor dSph.

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possible future extensions of this photometric and spectroscopic experiment.

2. SPECTROSCOPIC DATA SETS AND DATA REDUCTION

To collect low-resolution spectra of old (horizontal branch [HB]), \( V \sim 20.75 \) mag, intermediate (red clump [RC]; \( V \sim 20.5 \) mag), and young (blue plume [BP]; \( V \sim 22 \) mag) stellar tracers (see Fig. 1), we observed five fields with FORS2 at VLT (Appenzeller et al. 2008) in service mode.\(^{16}\) The program was almost completed by one of us (M. N.) during visitor-mode observations in 2004 December. A significant fraction of FORS2 spectra were collected using the 1400V grism, which covers H\( \beta \) (\( \lambda = 4861.34 \) Å) and Mg I Triplet (MgIT\(^{17} \)) lines, with a nominal resolution of 2100 (\( \lambda_c = 5200 \) Å). The slits used were 0.6” length, and typically 10” long, with a wavelength coverage of \( \approx 4560-5860 \) Å. The exposure times range from 22 minutes (short) to more than 1 hr (long).

During the 2004 December run, the 1028z grism was also used with a wavelength coverage \( \approx 7730-9480 \) Å. Its nominal resolution is 2560 (\( \lambda_c = 8600 \) Å) and covers the near-infrared (NIR) calcium triplet (CaT\(^{18} \)). Using the 1400V grism we secured 743 individual spectra of 356 stars with differential image motion monitor (DIMM) seeing ranging from 0.35″ to 1.27″ in 19.12 hr of exposure time, while using the 1028z grism, we secured 113 spectra of 73 stars with seeing 0.38″–0.59″ in 1 hr of exposure time. Most of the stars with 1028z spectra also have 1400V spectra. We ended up with a total sample of 359 stars either with single or multiple FORS2 spectra. In particular, the low-resolution (LR) data set collected with the grism 1400V includes 356 stars; among them, 11 have a single spectrum, while the others have a number of spectra ranging from 2 to 12. The LR data set collected with the grism 1028z includes 73 stars; among them, 39 have single spectra, while the others have a number of spectra ranging from two to three. In total, the LR data set includes 359 stars, of which 11 have a single spectrum and the others have a number of spectra ranging from 2 to 14. Table 1 gives the log of the observations of the spectra collected with FORS2. Listed, from left to right, are the date, the pointing, the coordinates, the grism, the exposure time, and the seeing.

These data were complemented by archival medium- and high-resolution spectra of red giant (RG) and RC stars collected with GIRAFFE (Pasquini et al. 2002) at VLT\(^{19} \) and covering the entire body of the galaxy (see Fig. 1). The GIRAFFE spectra (Koch et al. 2006) were collected using the grating LR8 with a nominal resolution of 6500 centered on the NIR CaT. This data set (hereafter GMR03) includes 10,394 individual fiber spectra of 1070 stars and were secured with a total exposure time of 95.4 hr with a seeing condition of 0.4″–2.2″. Note that this data set was also adopted by Muñoz et al. (2006) in their analysis of Carina radial velocity distribution.

We also use the high-resolution spectra collected with the gratings: HR10 (\( R = 19,800; 5339 \lesssim \lambda \lesssim 5619 \) Å), HR13 (\( R = 22,500; 6120 \lesssim \lambda \lesssim 6405 \) Å), and HR14A (\( R = 17,740, 6308 \lesssim \lambda \lesssim 6701 \) Å). This data set (hereafter GHR) includes 2002 individual spectra of 98 stars, secured with a total exposure time of 24 hr and seeing of 0.57″–3.20″. The preceding data were complemented by a new set of GIRAFFE spectra collected with LR8 and including 8092 individual spectra of 959 stars. This data set (hereafter GMR08) was secured with 48.4 hr of exposure time and seeing of 0.54″–1.94″. We ended up with a total sample of 1931 stars with either single or multiple GIRAFFE spectra.

The GMR ((GMR03 + GMR08), weighted mean) data set includes 1887 stars, and only one star has a single spectrum, while the others have multiple spectra ranging from 2 to 35. The HR data set includes 98 stars, and all of them have multiple spectra ranging from seven to nine. Moreover, the GHR and the GMR03 spectra have 38 objects in common, while 125 objects are in common between GHR plus GMR03 and GMR08 spectra. The LR resolution data set has 145 objects in common with the GIRAFFE (GHR plus GMR) data set. The objects in common were adopted to calibrate GMR and LR spectra.

Data plotted in the top panels of Fig. 1 show the spatial coverage of the different data sets, and the bottom panels show the location of the targets in the \( V \) and \( B - I \) color-magnitude diagram (CMD; Bono et al. 2010). The numbers in parentheses give the number of stars for which we estimated the radial velocity and the total number of stars for which we collected at least one spectrum.

Data plotted in the bottom panels of Fig. 1 show several interesting features:

**Stellar tracers.**—The FORS2 spectra cover old, intermediate, and young stellar tracers (\( V \lesssim 22 \), 359 RVs).

**Statistics.**—Although the exposure time of the GIRAFFE spectra is a factor of 4–8 longer than the FORS2 spectra, their limiting magnitude is \( V \lesssim 20.6 \) (fibers vs. slits). This means that this data set does not cover young (BP) and truly old (HB) tracers. The RGs are, indeed, a mix of old and intermediate-age stellar tracers. However, the sample of RV measurements based on GIRAFFE spectra (1985 vs. 359) is more than a factor of 5 larger than FORS2 (field of view of \( \approx 25 \) vs. \( \approx 5′ \) squared).

**Multiplicity.**—A significant fraction of our targets have multiple spectra. This means that we can provide robust estimates of intrinsic errors.

**Overlap of spectra resolution.**—We can constrain the occurrence of possible systematic errors in RV measurements using spectra collected with different instruments and different gratings/grisms.
Fig. 1.—Top left: Spatial distribution of different spectroscopic data sets. The plus signs and the triangles mark targets with medium- (LR03 and LR08) and high-resolution (HR10, HR13, and HR14A) GIRAFFE spectra. The background (gray dots) is a composite Carina reference image based on randomly selected subsamples of stars from the Bono et al. (2010) photometric catalog. The dashed ellipse shows the inclination ($i = 65^\circ$) and the radial extent (tidal radius, $r_t = 28.8 \pm 3.6'$; Mateo 1998). Top right: Same as the top left, but for targets with low-resolution FORS2 spectra (diamonds). The asterisks mark the positions of the targets with high-resolution spectra collected with MMFS at Magellan by W07. Bottom left: $V$ and $B - I$ CMD of spectroscopic targets. The symbols are the same as in the top left panel and show the position in the CMD of the targets with GIRAFFE spectra. The numbers in parentheses refer to the number of stars with measured radial velocities and to the total number of stars with at least one observed spectrum. Bottom right: Same as the bottom left, but for targets with FORS2 and MMFS spectra. See the electronic edition of the PASP for a color version of this figure.
We reduced our FORS2 data using standard IRAF\textsuperscript{20} tasks, using daytime-associated bias frames, flat field, and calibration lamps. The wavelength calibration was carried out using the task {	tidentify}, daytime lamp frames, and a set of 8–12 calibration lines. We typically end up with an accuracy better than \(\approx 0.02–0.04\, \text{Å}\) and \(\leq 2.5\, \text{km s}^{-1}\). However, we found systematic shifts up to 1 pixel, corresponding to \(\approx 30\, \text{km s}^{-1}\) in radial velocity, for the two adopted grisms. To overcome the problem, we adopted the approach suggested by Kelson (2003). A model of the sky for each slit was created and subtracted from the reduced 2D spectrum. This model was also used to estimate and correct the systematic shifts from the daytime wavelength calibration lamps, using night-sky lines, mainly O I 5577.34 Å for the 1400V grism and a set of isolated skylines for the 1028z grism.

Yet another potential source of systematics in RV estimates was the centroiding (see, e.g., Tolstoy et al. 2001). Thanks to the special care in the mask design, based on previous FORS2 pre-imaging of selected Carina fields, and to the very precise tracking of VLT, this effect was limited at the level of 5–10 km s\(^{-1}\). To estimate the shift between the centroid of the star and the slit, we adopted the slit image that is acquired just before the spectrum. For each target, the centroid of the slit is evaluated by collapsing the slit along its width and neglecting the pixels belonging to the star. The collapsed slit is then subtracted from the pixels belonging to the star, whose centroid is then estimated with the task {	t center} in IRAF. It is worth mentioning that we also acquired slit images for a couple of masks soon after we collected the spectra. These slit images were adopted to confirm that our estimates of the centroids were minimally affected by tracking problems. The details of the prereduction and reduction of FORS2 spectra will be discussed in a forthcoming article (Nonino et al. 2011, in preparation).

The GIRAFFE raw data were retrieved from the ESO Archive. These data were reduced using IRAF. After the standard bias and flat corrections, the spectra were extracted using the traces from flats. Wavelength calibration is based on daily calibration lamps: the formal solution rms was <0.02 Å, corresponding to a systematic <1 km s\(^{-1}\) in RV. After the fiberspectra assignment, performed using the table associated with the raw data, all the spectra were visually inspected in order to remove bad spectra. A master sky for each exposure was

\textsuperscript{20}IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.
created from sky fibers, after the spectra had been cross-correlated to remove shifts and scaled using the intensity of selected night-sky lines in the wavelength range of 8250–8750 Å, to better match the region with stellar absorption lines. Subsequently, the master skies were cross-correlated with the targets spectra, scaled in intensity, and subtracted. The cross correlation among the stacked skies showed that the relative accuracy in wavelength calibration could give systematics of the order of 500 m s\(^{-1}\) in RV. Finally, we co-added all the individual spectra of the same target using the task *scombine*, after correcting for the barycentric motion.

The GIRAFFE spectra taken with the high-resolution grisms were reduced following a similar approach.

### 3. RADIAL VELOCITY MEASUREMENTS

To measure the RV of GMR and GHR spectra, we fit individual spectral lines in the co-added spectra. The RV of the LR spectra was measured on individual spectra, and the final RV was estimated as a weighted mean among the multiple spectra, when available, of the same object. All spectra were normalized to the continuum and the fit to single or multiple lines was performed using either a Gaussian or a Moffat function. The interactive code developed to perform the RV measurements will be described in a forthcoming article (Fabrizio et al. 2011, in preparation). By visual inspection, we selected several spectral lines for each data set (namely, H\(\beta\) and Mg I triplet for FORS2/1400V spectra; see the spectra plotted in the panels of Fig. 2d), the CaT for FORS2/1028z and GMR spectra (see the spectra plotted in Figs. 2b and 2c), and several strong Fe I lines for GHR spectra (see the spectra plotted in the panels of Fig. 2a). The quoted lines were typically fit with a Moffat function with \(\beta = 2\), to properly account for the contribution of the wings in the line fit.

This preliminary estimate of the RV was validated by visual inspection, and when judged to be satisfactory, it was adopted to perform an automatic estimate of the RV using a large set of iron and heavy-element lines. We adopted the line list for iron, \(\alpha\), and heavy elements recently provided by Romaniello et al. (2008) and by Pedicelli et al. (2010). The lines were selected according to the wavelength range and the spectral resolution of the different instruments and grisms. We ended up with a subsample of 30–70 lines for LR (grisms: 1028z, 1400V) spectra, with \(\approx 35\) lines for GMR spectra and with \(\approx 90\) lines for the GHR spectra. For each spectrum we estimate the RVs as a weighted mean of the lines with the highest signal-to-noise ratio (S/N). On average, this means 1–2 dozen lines for LR and GMR spectra and approximately 40 lines for the GHR spectra.

The error in the RV based on individual lines was assumed to be equal to the sigma of the fitting function.\(^{21}\) The RV of the entire spectrum was estimated as a weighted mean over the different fitted lines. The error in the radial velocity of co-added spectra (GMR and GHR) is the error of the weighted mean. The RV of the LR spectra was estimated following the same approach for stars with a single spectrum and as a weighted mean among the different mean measurements for the stars with multiple spectra.

Data plotted in Fig. 3 show that the intrinsic error of the RV measurements based on GHR spectra is, on average (biweight mean), \(\approx 1\) km s\(^{-1}\) (Figs. 3a, 3b, and 3c) and becomes of the order of \(5\) km s\(^{-1}\) for GMR03 spectra (Fig. 3d). Note that the intrinsic error of GMR08 spectra is larger and \(\approx 6\) km s\(^{-1}\). The difference is mainly due to the S/N, since the latter sample is, on average, 0.5 mag fainter, and its total exposure time is almost a factor of 2 shorter.

The intrinsic error of the RVs based on LR spectra is slightly larger and ranges from \(\approx 11\) km s\(^{-1}\) (grism 1028z, Fig. 3f) to \(\approx 10\) km s\(^{-1}\) (grism 1400V, Fig. 3g). The difference between the FORS2 data sets is expected. The first two lines of the CaT, at the typical RV of Carina stars, are contaminated by atmospheric OH emission lines. Therefore, the RV measurements based either on low-resolution spectra or on medium-resolution, low S/N spectra might be affected by slightly larger uncertainties in case the sky lines are not precisely subtracted. The reader interested in a more detailed discussion is referred to Nonino et al. (2007) and to W07. Moreover, the 1028z spectra have a number of repeats that are approximately a factor of 2 smaller than the number of repeats of the 1400V spectra (see § 2).

The fraction of objects for which we measured the RV is quite high; indeed, the fraction ranges from \(\approx 98\%\) for GHR and GMR03 spectra to \(\approx 90\%\) for LR (1400V) spectra and to more than \(\approx 80\%\) for GMR08 and LR (1028z) spectra. The stars with no RV measurements typically have either low S/N, noisy spectra (see § 2), or missing identification in the photometric catalog. Eventually, the RV was measured in 2165 out of the 2323 GIRAFFE spectra (\(\approx 93\%\)) and in 381 out of the 429 (\(\approx 89\%\)) FORS2 spectra (see Fig. 3). In total, the RV was measured in 1812 stars using GIRAFFE spectra and in 324 stars using FORS2 spectra (see Fig. 1). The final catalog includes RV measurements for 1979 stars.

To validate the approach adopted to measure the RVs, we decided to use real spectra. This approach has several indisputable advantages when compared with artificially generated spectra, since the different spectroscopic data sets we are dealing with partially overlap. We can constrain the precision of RVs based on HR spectra, since for each object we have three spectra with similar spectral resolutions and covering three different

\(^{21}\)In particular, we assume \(\epsilon_\text{RV} = \sigma_\lambda / \lambda\) km s\(^{-1}\).
wavelength regions (HR10, HR13, and HR14A; see § 2). Figures 4a and 4b show the internal comparison of our RV measurements as a function of the V-magnitude. The difference and the RV dispersion of the entire sample and of the candidate Carina stars (180 ≤ RV ≤ 260 km s⁻¹/C0; see § 4) are minimal.

We can also constrain the precision of the GMR03 spectra, since this subsample has more than three dozen stars in common with the GHR spectra. Figure 4c shows the difference between the RVs based on the (GHR) and those based on the GMR03 spectra (weighted mean). The mean (biweight) of the difference is 0.1 km s⁻¹ for the candidate Carina stars (26) and smaller than 1 km s⁻¹ for the entire sample (35). The GMR03 RVs were corrected for the difference.

Figure 4d shows the difference between the RVs based on (GHR) plus GMR03 (weighted mean) spectra and RVs based on GMR08 spectra. We have more than 120 stars in common and all of them are candidate Carina stars. The difference is smaller than −3 km s⁻¹ and the σ ≈ 8 km s⁻¹. The difference with the RVs based on GMR03 spectra is mainly due to the fact that the GMR08 spectra have a lower S/N and are, on average, fainter. The GMR08 RVs were corrected for the difference.
FIG. 5a shows the difference between the RVs based on the GIRAFFE (\( \langle GHR + GMR \rangle \), weighted mean) spectra and on the FORS2-1400V spectra (internal comparison). The mean difference (biweight) over the entire sample of stars in common (\( 140 \leq RV \leq 144 \)) is \( -3.6 \) km s\(^{-1} \) and smaller (\( -2.7 \) km s\(^{-1} \)) for the candidate Carina stars (\( 122 \leq RV \leq 125 \)), while the dispersion is of the order of 12 km s\(^{-1} \). The RVs based on FORS2-1400V were accordingly corrected. Figure 5b shows the difference between the RVs based on GIRAFFE and on FORS2 (\( \langle GHR + GMR \rangle + 1400V \), weighted mean) and RVs based on FORS2-1028z spectra. The difference in RV and dispersion are similar to the RVs based on the 1400V grism and they were also corrected. Data plotted in Fig. 5c show the difference between the RVs based on the GIRAFFE (\( \langle GHR + GMR \rangle \), weighted mean) and on the FORS2 (\( (1028z + 1400V) \), weighted mean) spectra after the corrections have been applied. The two samples have more than 140 stars in common and the difference is \( \sim 1 \) km s\(^{-1} \) for the entire sample and vanishing for candidate Carina stars.

Figure 6 shows the comparison between our RVs based on GIRAFFE spectra—\( \langle GHR + GMR \rangle \)—with different RV measurements available in the literature (external comparison). Figure 6a shows the comparison with the RVs based on echelle (“\( +2D-\text{FRUTTI} \)” at the Las Campanas Observatory 2.5 m du Pont Telescope; Mateo et al. 1993 [crosses]) and on multifiber spectra (Hydra at the Cerro Tololo Inter-American Observatory 4 m Blanco Telescope; Majewski et al. 2005 [filled circles]). The difference with the Mateo et al. (1993) sample is quite small, both in the mean (biweight) and in the velocity dispersion. The difference with the Majewski et al. (2005) sample is larger (\( \mu = -3.6 \) and \( \sigma = 9.4 \) km s\(^{-1} \)), and it is caused by the lower S/N of the spectra they adopted to estimate the RVs. Figure 6b shows the comparison between our RVs and RVs based either on high-resolution spectra collected with MIKE at
Magellan (crosses) or on medium-resolution spectra collected with GIRAFFE at VLT (filled circles) provided by Muñoz et al. (2006). The MIKE spectra were collected in slit mode with a resolution of $\sim 19,000$ (red arm) and the RV measurements are based on the NIR calcium triplet. The GIRAFFE spectra adopted by Muñoz et al. (2006) are the same GMR03 spectra we also included in this investigation. The difference with RVs based on MIKE spectra is minimal ($<1$ km s$^{-1}$), but the sample of stars in common is limited (14). The difference with the RVs based on GIRAFFE spectra is larger $\sim 3$ km s$^{-1}$ if we account for the entire sample of stars in common (850/889 stars, $\sigma = 5.7$ km s$^{-1}$) and for candidate Carina stars (353/373 stars, $\sigma = 6.7$ km s$^{-1}$). The reasons for this difference are not clear, apart from the fact that we adopted a different approach to pre-reduce the spectra and to measure the RVs.

Figure 6e shows the comparison with the RVs provided by W07. The RV measurements provided by W07 are based on high-resolution (20,000–25,000) spectra collected with MMFS at Magellan. These spectra sample the region across the magnesium triplet (5140–5180 Å). We have more than 550 candidate Carina stars in common, and the data plotted in this figure show that the two samples of RV measurements agree, within the errors, quite well. The difference is $-2$ km s$^{-1}$ (biweight), and the RV dispersion is smaller than 10 km s$^{-1}$. The same outcome applies to the RVs based on FORS2 spectra (see Fig. 6d).

We have more than 80 candidate Carina stars in common, and the difference is $\approx 1$ km s$^{-1}$, while the dispersion $\sigma \sim 11$ km s$^{-1}$.

We ended up with a data set including RV measurements for 1979 stars. For the objects with RVs based on both GIRAFFE and FORS2 spectra we performed a weighted mean, with the inverse square of the RV standard errors of the GMR, GHR, and LR spectra used as weights. As a final test, we compared our entire RV data set with the RVs provided by W07. Data plotted in Figure 6e further support the agreement between the two different sets of RV measurements; indeed, the difference ranges from less than...
1 km s$^{-1}$ for the entire sample (785 objects in common) to 2 km s$^{-1}$ for candidate Carina stars (574), while the dispersion is $\sim$9 km s$^{-1}$. The estimates of the mean difference are once again based on the biweight, and the number of neglected objects is smaller than 9%. The precision of both internal and external validations and, in particular, the good agreement with the RVs provided by W07 support the approach we adopted to measure the RVs.

The occurrence of binary stars among RG stars in dSph galaxies is a highly debated topic, and their impact on the RV dispersion ranges from a sizable (Queloz et al. 1995) to a small error (Olszewski et al. 1996; Hargreaves et al. 1996) when compared with the statistical error. However, in a recent investigation, by using a detailed statistical approach, Minor et al. (2010) found that dSph galaxies with RV dispersions ranging from 4 to 10 km s$^{-1}$ can be inflated by no more than 20%, due to the orbital motion of binary stars. It is worth noting that the current LR spectra cover a time interval of 3 yr (2004–2007) and the number of repeats ranges from 2 to 14 (only 11 stars with a single spectrum). The GMR spectra cover a time interval of 5 yr (2003–2008) and the number of repeats ranges from 2 to 35 (only one star with a single spectrum). The GHR spectra were collected in 2005, and the number of repeats ranges from seven to nine. The three different data sets have 145 objects in common. This means that a significant fraction of stars in our sample have spectra collected on a time interval of several years. Therefore, the current RV measurements are less prone to significant changes caused by binary stars. Moreover, the conclusions of the current investigation are minimally affected by a possible uncertainty of the order of 20% in the RV dispersion.

4. CARINA RADIAL VELOCITY DISTRIBUTION

The top panel of Figure 7 shows the RV distribution of the entire (GIRAFFE plus FORS2) data set. The well-defined primary peak located at RV $\sim$220 km s$^{-1}$ includes candidate Carina stars. The current radial velocity distribution is soundly supported by the radial velocity distribution recently provided by W07, but based on a smaller number of stars (see the middle panel of Fig. 7). This evidence and the minimal difference in the RV of the stars in common allowed us to merge the two RV catalogs. For the objects in common, weighted mean RVs were computed with the inverse square standard errors in the individual RV measurements used as weights. We ended up with a sample of 2629 stars (see the bottom panel of Fig. 7).

To further constrain the radial velocity distribution of the candidate Carina stars, we ran a Gaussian kernel on each star with a $\sigma$ equal to the RV uncertainty. The solid line plotted in the top panel of Figure 8 was computed by summing the individual Gaussians over the entire data set. We estimated the mean and the $\sigma$ by fitting the smoothed RV distribution with a Gaussian (dashed line). We found that the peak (220.4±0.1 km s$^{-1}$) in the RV distribution agrees quite well, within the errors, with similar estimates available in the literature (Koch et al. 2006; Muñoz et al. 2006; W07; Walker et al. 2009). The same outcome applies to the RV dispersion. Note that current estimate is based on a homogeneous sample of candidate Carina stars (180 $\leq$ RV $\leq$ 260 km s$^{-1}$, $\sim$4$\sigma$, 1208 stars) that is $\sim$55% larger than any previous sample of RVs.

The middle panel of Figure 8 shows the RV distribution of candidate Carina stars, but based on the RV measurements (780 candidate Carina stars) provided by W07. The two data sets provide, within the errors, very similar kinematic properties. Note that by merging the two data sets of RV measurements, we ended up with a sample of 1378 candidate Carina stars according to the quoted selection criterion. This sample of Carina RVs is $\sim$75% larger than any previous RV sample. To overcome subtle uncertainties in the estimate of the RV due to the possible presence of outliers in the final data set, we estimated the biweight mean of this sample and we found 220.9 ± 0.1 km s$^{-1}$. This estimate is based on 1369 stars, and this is
the sample that we will adopt in the following to estimate the RV distribution (solid line in the bottom panel of Fig. 8). The mean based on the Gaussian fit (dashed line) agrees quite well with the biweight mean. However, the tails of the smoothed RV distribution appear shallower than the tails of the fitting Gaussian. There is also evidence that the low-velocity tail might be shallower than the high-velocity one. To provide a more quantitative estimate, we estimated the $\chi^2$ of the two curves, and we found that the hypothesis that the smoothed distribution is Gaussian can be discarded at 99% confidence level.

To constrain the possible dependence of the Carina RV distribution on the different samples of RV measurements, we estimated the RV distribution of Carina stars using stars brighter than $V = 20.25$ mag (651). We fit the smoothed distribution with a Gaussian and we found a mean RV $\mu = 221.3 \pm 0.1$ km s$^{-1}$ and a $\sigma = 8.7 \pm 0.1$ km s$^{-1}$. We also performed a new test including only stars brighter than $V = 20.50$ mag (1050) and we found $\mu = 220.9 \pm 0.1$ km s$^{-1}$ and $\sigma = 8.9 \pm 0.1$ km s$^{-1}$. The quoted values agree quite well with the estimates based on the entire sample. The two RV distributions based on smaller samples of brighter objects also show asymmetric wings when compared with a Gaussian distribution.

To further determine the plausibility of the selection criterion to pinpoint candidate Carina stars, we plotted the kinematic candidates in the $V$ and $B - I$ (left panel of Fig. 9) and in the $V$ and $V - I$ (right panel of Fig. 9) CMDs. The symbols of the different spectroscopic data sets are the same as in Figure 1. Data plotted in this figure quite clearly show the accuracy of the kinematic selection; indeed, the bulk of the candidate Carina stars are located along the expected evolutionary sequences (RG, RC, HB, and BP). The new sample also provides a clear separation between RG and asymptotic giant branch (AGB; $V \sim 20$ and $B - I \sim 1.6$--1.9 mag) stars. Moreover, and even more importantly, data plotted in Figure 9 soundly support the results recently provided by Bono et al. (2010) concerning the metallicity distribution of Carina stars. The quoted authors found, using a color-color plane ($U - V$ and $B - I$) to select candidate Carina stars, that the spread in color of RG stars is significantly smaller than suggested by spectroscopic measurements. The very narrow distribution in color of candidate Carina RGs is further supported by the kinematic selection. There are a few RGs that attains colors that are redder than typical Carina RGs. They might be either variables stars or misidentified objects.

5. DISCUSSION

In order to determine whether the occurrence of asymmetries in the radial velocity distribution is caused by the presence of substructures in Carina, we investigated the change in the RV distribution as a function of the projected radial distance ($r$). To avoid spurious fluctuations in the mean RV, we ranked all the candidate Carina stars as a function of $r$ and estimated the mean (biweight) and the median using the first 200 objects. The $\rho$ of each subsample was estimated as the mean over the individual distances of the same 200 stars. We estimated the same quantities by moving one object in the ranked list until we accounted for the most distant 200 stars in our sample located inside the tidal radius ($r_t \sim 28.8 \pm 3.6'$; Mateo 1998). The error on the mean RV for individual bins is smaller than one-tenth of km s$^{-1}$. In order to provide robust constraints on the possible uncertainties introduced by the number of stars per bin and by the number of stepping stars, we performed a series of Monte Carlo simulations. The estimated mean dispersion is plotted as a shaded area across the RV mean in the bottom panel of Figure 10. Note that changes in the binning criteria can change the extent of the secondary features, but their relative positions are minimally affected.

The data plotted in the bottom panel of Figure 10 show that the mean RV attains a well-defined minimum—
RV $\sim 219.5 \pm 0.4$ km s$^{-1}$ for $\rho \sim 7^{\prime}$ (200 pc) from the galaxy center and increases up to $\sim 222.6 \pm 0.4$ km s$^{-1}$ for $\rho \sim 13^{\prime}$ (400 pc, 6$\sigma$ level) and a plateau of $\sim 221.4 \pm 0.4$ km s$^{-1}$ at larger distances (600 pc, 4$\sigma$ level). This evidence is robust, since the mean (biweight) and the median attain very similar values and intrinsic errors. To further constrain the change of the kinematic status across the body of the galaxy, we also estimated the radial velocity dispersion ($\sigma_{\text{RV}}$, solid line in the middle panel of Fig. 10). Note that to avoid subtle uncertainties caused by asymmetric radial velocity distributions, we also estimated for each bin the semi-interquartile range (SIQ; dashed-dotted line in the middle panel of Fig. 10). Data plotted in the middle panel of Figure 10 show that both the $\sigma_{\text{RV}}$ and the SIQ show very similar radial trends, but the former dispersion parameter attains values that are $\approx 50\%$ larger. The difference indicates that the velocity dispersion should be cautiously treated in dealing with asymmetric RV distributions (Strigari et al. 2010). We also found that the dispersion is larger in the innermost regions, decreases in coincidence of the minimum in the mean RV, and (after a mild increase) attains an almost constant value in the outermost galaxy regions. Note that the shaded areas across the RV dispersion and the SIQ display the intrinsic error on individual bins according to Monte Carlo simulations. There is a hint that the dispersion could increase in the outskirts, thus supporting the finding by W07. However, more data are required to constrain the change in the projected radial velocity dispersion across the tidal radius.

Interestingly enough, the kurtosis plotted in the top panel of Figure 10 shows a mirror trend, compared with the projected radial velocity dispersion and with the mean RV. This suggests that in the Carina regions where the mean RV is smaller, the dispersion is also smaller and the radial velocity distribution is more centrally peaked. The opposite trend characterizes the regions in which the mean RV is larger. The outermost regions show flat trends in the quoted parameters. Evidence of kinematic and photometric substructures have already been found in several LG dwarfs (Monelli et al. 2003; Kleyna et al. 2004; Walker et al. 2009). In particular, Battaglia et al. (2008) found evidence of a radial velocity gradient in Sculptor that they interpreted as an indication of internal rotation. The current data indicate that the RV distribution across the body of Carina shows a radial gradient and possible evidence of rotation. However, the RV distribution is far from being uniform, possibly suggestive of the occurrence of well-defined spatial substructures.

To clarify the approach we adopted to estimate the weighted mean between the GIRAFFE and the FORS2 spectra, Figure 11 shows the same data as Figure 10, but the parameters plotted in the left panels are only based on the GIRAFFE spectra, while the right panels are only based on the FORS2 spectra. Figure 9 shows the $V$ and $B-I$ (left) and $V$ and $V-I$ (right) CMDs of candidate Carina stars selected on the basis of the radial velocity ($180 \leq \text{RV} \leq 260$ km s$^{-1}$, $\sim 4\sigma$). The symbols of the different spectroscopic data sets are the same as in Fig. 1. See the electronic edition of the PASP for a color version of this figure.

\footnote{Projected distances on Carina were estimated assuming a true distance modulus of 20.10 $\pm$ 0.12 mag (Dall’Ora et al. 2003; Pietrzyński et al. 2009) together with the inclination angle and the axes ratio (Mateo 1998).}
those plotted in the right panels are based on the weighted mean between GIRAFFE and FORS2 spectra. The trends of the plotted parameters as a function of the radial distance are the same. They are only slightly more noisy, due to the decrease in the sample size.

To further constrain the spatial extent of possible kinematic substructures, we investigated the same parameters, but they were estimated as a function of right ascension (\(\alpha\)) and declination (\(\delta\)). Data plotted in Figure 12c show the marginal of the radial velocity along the right ascension axis. We found that the mean RV decreases from \(221.9 \pm 0.4 \text{ km s}^{-1}\) when moving from east (~750 pc from the galaxy center) to west (~1060 pc from the center). The difference is equal to ~2.1 \text{ km s}^{-1} (4\sigma\ level) and increases by almost a factor of 2 (~4.5 \text{ km s}^{-1}, 9\sigma\ level) if we take into account the absolute maximum (~223.3 \pm 0.4 \text{ km s}^{-1}, ~320 pc from the center) and the absolute minimum (~218.8 \pm 0.4 \text{ km s}^{-1}, ~280 pc from the center) of the mean RV.

It is worth mentioning that the increase in the mean velocity in both the external regions and outside the Galactic center appears to be in the same direction of the Carina proper motion (Fig. 12d [long arrow]; Pi
tek et al. 2003; Walker et al. 2009). However, the uncertainty affecting current estimates of the Carina proper motion is quite large (short arrows). Note that the increase in the mean RV cannot be in the direction of the Galactic center (heavy dashed line). Moreover, and even more importantly, in a region of 200 pc across the Galactic center we found that the RV shows a secondary maximum that is strongly correlated with the velocity dispersion (Fig. 12b) and anticorrelated with the kurtosis (Fig. 12a).
The change in the mean RV along the declination axis is similar to the change in the right ascension axis. The mean RV decreases when moving from north to south (Fig. 12e), from 221.9 ± 0.4 (≈630 pc from the center) to 219.8 ± 0.4 km s⁻¹ (≈−650 pc from the center). The difference is at least at 4σ level, while the difference between maximum and minimum (222.3 ± 0.4 km s⁻¹ and ≈238 pc; 219.1 ± 0.4 km s⁻¹, and ≈41 pc) is more than 3 km s⁻¹ (∼7σ level). The absolute minimum is located close to the center in the IV quadrant and associated with the absolute minimum in right ascension. The

Fig. 12—(a–c) Same as Fig. 10, but as a function of right ascension (bottom axis) and distance in parsecs (top axis). (d) Sky distribution of the different spectroscopic data sets. Symbols and colors are the same as in Fig. 1. The long arrow shows the Carina proper motion according to Piatek et al. (2003) and Walker et al. (2009), and the short arrows display current uncertainties. The dashed line shows the direction of the Galactic center according to an observer located in the center of Carina. The dotted and the dash-dotted thin lines display the secondary features identified in Fig. 10 and in Fig. 13. (e–g) Same as Fig. 10, but as a function of declination (left axis) and distance in parsecs (right axis). See the electronic edition of the PASP for a color version of this figure.
radial trend outside this region shows the quoted secondary maximum and a constant smooth trend at larger distances from the center. The dispersion (Fig. 12f) and the kurtosis (Fig. 12g) are once again correlated/anticorrelated with the behavior of the mean RV. It is noteworthy that the quoted findings are minimally affected by the FORS2, GIRAFFE, and W07 data sets adopted to estimate the RVs of candidate Carina stars.

To explain the occurrence of a kinematically cold population close to the center of Sextans, Kleyna et al. (2004) suggested, using also photometric radial distributions, that it could be the aftermath of a merging with a globular cluster. To further obtain the nature of the preceding kinematic peculiarities, we investigated the Carina radial distributions using the same approach adopted by Monelli et al. (2003) and the recent photometric catalog by Bono et al. (2010). The isocontour levels (Fig. 13b) and the marginal (Fig. 13a) along the right ascension axis plotted in Figure 13 show that the radial distribution is flat-topped across the Carina center. Note that the bright field star HD48652 ($V = 9.14$ mag) minimally affects this feature, since it is located at 2.77′ from the Carina center. On the other hand, the radial distribution is symmetric and with a well-defined peak along the declination axis. We drew three vertical lines across the center and the two secondary maxima. Interestingly enough, the same lines include the secondary maximum detected along the right ascension axis in the mean RV. These empirical evidence indicates that the regions across the galaxy center might form a substructure probably reminiscent of a transition between a bulgelike and/or a disklike substructure.

Fig. 13.—(a) Normalized radial profile along the right ascension axis of the Carina photometric catalog (Bono et al. 2010). (b) Sky distribution of the photometric catalog. The contours display isodensity levels ranging from 5 to 95% with steps of ~11%. The two thin dash-dotted lines mark the two secondary peaks, and the thin dotted line marks the center of the galaxy. The star shows the position of the bright field star HD48652 ($V = 9.14$ mag). (c) Same as (a), but along the declination axis. See the electronic edition of the PASP for a color version of this figure.
6. SUMMARY AND CONCLUSIONS

We presented new radial velocity measurements of Carina dSph stars based on low-resolution spectra (R ~ 2500) collected with the FORS2 multiobject slit spectrograph at the VLT. The key advantage of the current sample is that for the first time we collected spectra of old (HB, V ~ 20.75 mag), intermediate (RC, V ~ 20.5 mag), and young (BP, V ~ 22 mag) stellar tracers. The bulk of the spectra were collected using the 1400V grism covering the Hβ and the Mg I triplet. For a fraction (20%) of these targets we also collected spectra using the 1028z grism centered on the NIR CaT. We secured 856 spectra of 359 stars with ~20 hr of exposure time.

The preceding data were complemented by medium-(R ~ 6500) and high-resolution (R ~ 18,000–23,000) spectra collected with the GIRAFFE multiobject fiber spectrograph at the VLT. The targets of this spectroscopic data set are RG and RC stars (V ≤ 20.6 mag). The medium-resolution spectra are centered on the NIR CaT, and the high-resolution spectra cover a broad wavelength region (5300 ≤ λ ≤ 6700 Å). In total, this data set includes 20,488 spectra of 1931 stars collected with ~168 hr of exposure time.

The radial velocity distribution based on the preceding spectra agrees quite well with similar data available in the literature. By assuming that candidate Carina stars are objects with RVs ranging from 180 to 260 km s\(^{-1}\) (4σ), we ended up with a sample of ~1,210 stars. By fitting these objects with a Gaussian, we found a mean RV—(RV) = 220.4 ± 0.1 km s\(^{-1}\)—and a velocity dispersion—σ = 11.7 ± 0.1 km s\(^{-1}\)—that agree, within the errors, with similar measurements provided by W07.

The preceding kinematic selection was firmly supported by the position of the candidate Carina stars along the expected evolutionary sequences (RG, AGB, RC, HB, and BP) of the CMD. The narrow distribution in color of the candidate Carina RG stars confirms the recent findings by Bono et al. (2010), based on photometric selection criteria, that the spread in metallicity of Carina stars is smaller than suggested by spectroscopic measurements (Koch et al. 2006).

The current catalog of Carina RV measurements was complemented by similar measurements recently provided by W07. The stars in common were averaged and we ended up with a sample of ~1370 RV measurements of candidate Carina stars that is ~75% larger than any previous Carina RV sample. We found that the hypothesis that the Carina RV distribution is Gaussian can be discarded at the 99% confidence level. A more detailed investigation of the RV across the body of the galaxy indicates that the mean RV changes from ~220 km s\(^{-1}\) at a distance of 7' (~200 pc) from the center to ~223 km s\(^{-1}\) at 13' (6σ level) and attains an almost constant value of ~221–222 km s\(^{-1}\) at larger distances (3–5σ level). Moreover, we found that in the Carina regions where the mean RV is smaller, the dispersion also attains smaller values and the RV distribution is more centrally peaked (larger kurtosis).

The mean RV at large distances from the Galactic center decreases by more than 2 km s\(^{-1}\) (4σ level) when moving from east to west. The difference increases by almost a factor of 2 (~4 km s\(^{-1}\), 9σ level) if we account for the difference between the absolute maximum and the absolute minimum of the mean RV. The change in the mean RV is similar when moving from north to south. The difference increases from more than 2 km s\(^{-1}\) (4σ level) in the outermost regions to more than 3 km s\(^{-1}\) (7σ level) across the galaxy center. The RV gradient appears to be in the direction of the Carina proper motion. However, this parameter is still too affected by large uncertainties to reach a firm conclusion.

Furthermore, there is also evidence of a secondary maximum in RV across the Carina center (|D| ≤ 200 pc). The occurrence of a secondary feature across the Carina center is also supported by the flat-topped radial distribution (photometric catalog) along the right ascension axis. In particular, the two secondary maxima cover the same galaxy regions of the secondary maximum in the RV radial variation. These results, once independently confirmed, are probably reminiscent of a substructure with transition properties.

Evidence of substructures in dwarf galaxies dates back to Jerjen et al. (2000), who detected a spiral structure in the Virgo dE IC3328 using deep ground-based R-band images. This discovery was soundly confirmed by Lisker et al. (2007), who found that a significant fraction of bright (M_B < 16) early-type dwarfs in Virgo are characterized by disklike features (spiral arms, bars). Furthermore, they also found that the properties of dE in clusters are also correlated with environmental density. This evidence supports numerical simulations of galaxy harassment indicating that late-type galaxies undergo a significant transformation when accreted in a cluster (Mastropietro et al. 2005; Moore et al. 1996, 1998). It was also suggested that dEs with disklike features might be the low-luminosity tail of normal disk galaxies (Lisker et al. 2006). This working hypothesis was recently supported by new and accurate kinematic data of IC3328 suggesting that the observed velocity dispersion is the aftermath of two distinct substructures, namely a thin stellar disk and a dynamically hot component. However, no robust conclusion could be reached, due to limited radial coverage (Lisker & Fuchs 2009).

Irregular kinematic properties have also been found in the recently discovered ultrafaint dwarf galaxy—Willman I—by Willman et al. (2010) using both photometric and spectroscopic data. They found that galaxy stars located in innermost regions show a radial velocity offset of 8 km s\(^{-1}\) when compared with the outermost ones. Moreover, they also found initial hints of asymmetries in the radial velocity distribution, but their spectroscopic sample is too small to reach firm conclusions.

If the complex kinematic properties of Carina will be supported by future more accurate and deep data sets, then the occurrence of disklike features might be considered a widespread property of gas-poor dwarf galaxies. A detailed comparison
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