2D numerical investigations derived from a 3D dragonfly wing captured with a high-resolution micro-CT

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Abstract.
BACKGROUND: Due to their corrugated profile, dragonfly wings have special aerodynamic characteristics during flying and gliding. OBJECTIVE: The aim of this study was to create a realistic 3D model of a dragonfly wing captured with a high-resolution micro-CT. To represent geometry changes in span and chord length and their aerodynamic effects, numerical investigations are carried out at different wing positions. METHODS: The forewing of a Camacinia gigantea was captured using a micro-CT. After the wing was adapted an error-free 3D model resulted. The wing was cut every 5 mm and 2D numerical analyses were conducted in Fluent\textsuperscript{R} 2020 R2 (ANSYS, Inc., Canonsburg, PA, USA). RESULTS: The highest lift coefficient, as well as the highest lift-to-drag ratio, resulted at 0 mm and an angle of attack (AOA) of 5°. At AOAs of 10° or 15°, the flow around the wing stalled and a Kármán vortex street behind the wing becomes visible. CONCLUSIONS: The velocity is higher on the upper side of the wing compared to the lower side. The pressure acts vice versa. Due to the recirculation zones that are formed in valleys of the corrugation pattern the wing resembles the form of an airfoil.

Keywords: Dragonfly wing, gliding flight, insect flight, computational fluid dynamics, aerodynamics

1. Background

Considering aerodynamic characteristics, the dragonfly with its special wing structure has one of the most interesting ways of flying and gliding. Its corrugated profile leads to special aerodynamic characteristics and manoeuvrability [1]. In previous studies, structural analyses were conducted to investigate the mechanical properties of the wing [1–6]. Next to the structural influence of the corrugations, the aerodynamic effects of the corrugations also were investigated during gliding flight [7–12]. Therefore, the aerodynamic performance of the pleated dragonfly wing was compared to streamlined airfoils [11,12].
Furthermore, it was additionally compared to flat plates [9,10]. Lift and drag coefficients were calculated under different flow conditions [9–12]. Okamoto et al. compared two- and three-dimensional (2D and 3D) models during gliding flight in a wind tunnel [13]. The hindwing used by Okamoto et al. was reconstructed in Chen and Skote’s study [14]. Only the first corrugations were considered. The further wing structure in the direction of the chord length was neglected and assumed to be a membrane [14]. It was found that only a 3D model of the whole wing is able to capture the occurring changes in spanwise direction [14]. Next to gliding flight, flapping flight has also been investigated [15–18]. Song et al. conducted numerical studies on a 3D forewing model neglecting small corrugated structures and small variations in the vein frame [15]. Dragonflies are a great role-model for bionic micro-air-vehicles (MAVs) [14,19,20]. MAVs and Biomimetic micro-air-vehicles (BMAs) are very light and small features, which can imitate the movements of the wings of birds, insects, and bats [20]. To fully capture all aerodynamic effects during gliding flight, it is important not only to reconstruct and approximate a wing but to represent the actual geometry with its corrugations and real form of the vein structure. Therefore, the aim of this study is to create a model of a real dragonfly wing to perform Computational Fluid Dynamics (CFD) investigations.

2. Objective

This work aimed to create a 3D model of a dragonfly wing as it occurs in nature. The 3D wing was cut at several locations to represent geometry changes and their aerodynamic effects in span length and chord length. With these 2D cross-sections, numerical analyses were conducted.

3. Methods

The forewing of a Camacinia gigantea was captured using a micro-CT. Due to the limited resolution of the micro-CT with 7 µm and the small dimensions of the dragonfly wing, only the vein structure of the wing could be captured. Therefore, the membrane needed to be created and adapted manually. By placing surfaces in the centerlines of the vein structure, the membrane was created. The membrane was adapted and thickened in Netfabb® Premium 2020 (Autodesk, Inc., San Rafael, CA, USA) and Meshmixer (Autodesk, Inc., San Rafael, CA, USA). Afterward, the membrane and the vein structure were connected using a Boolean union. Due to the Boolean operation, some intersections occurred in the areas where the membrane and the vein structure were united. These intersections and errors were removed using an adapted pymesh script [21]. After the pymesh script was applied to the geometry, most of the errors were removed. Although a theoretically valid STL-geometry could be achieved it was not possible to create a surface mesh in CENTAUR™ 14.5.0.2 (CentaurSoft, Austin, TX, USA), due to local STL inconsistencies. These problematic areas were manually corrected in CADfix (ITI, a Wipro company, Milford, OH, USA).

The 3D wing was cut at several locations to represent geometry changes in chord length and span length as well as their aerodynamic effects. As an initial study, only 2D results are presented within this short communication. The origin of the coordinate system is located at half of the span length. Therefore, cross-sections were created at 5 mm intervals, starting at the wing root at −20 mm to the wing tip at 15 mm.

The cross-sections were converted to splines for meshing in CENTAUR™. Angles of attack (AOA) of 0°, 5°, 10°, and 15° were set for the wing cross-sections. The sections with AOAs of 0° and 5° were analyzed using steady conditions, whereas the sections with AOAs of 10° and 15° were investigated
transiently with a time step size of $5 \times 10^{-5}$ s to resolve unsteady vortex motion. The calculations were performed incompressible and laminar. The wing is surrounded by a bounding box with a velocity inlet of 3 m/s and a pressure outlet with a gauge pressure of 0 Pa. The upper and lower boundaries of the bounding box were defined as euler wall. The solution was initialized using hybrid initialization.

For pressure-velocity coupling, the SIMPLEC scheme was used for the steady and transient calculations. To calculate the gradient, the Least Squares Cell Based method was chosen, for momentum second order upwind and second order implicit for the transient formulation. A grid convergence study was performed leading to an accuracy of one drag count. Convergence within the range of $1 \times 10^{-5}$ was reached. The results and images were plotted in Tecplot® 360 EX 2020 R2 (Tecplot, Inc., Bellevue, WA, USA). In addition, the resulting pressure coefficients were plotted in Tecplot® and qt grace.

4. Results

Figure 1a and b show the resulting STL-geometry. Figure 1c shows a detail of the resulting mesh at location 5 mm. Steady analyses were conducted for the wing cross-sections at AOAs of 0° and 5°. All lift coefficients at an AOA of 5° are higher than the lift coefficients of the same cross-section at 0° (see Fig. 2a and b). Figure 3 shows the velocity magnitude contour of the wing cross-section at 0 mm and an AOA of 5°. The highest lift coefficient, as well as the highest lift-to-drag ratio, resulted there. Comparing the results for lift and drag coefficients with the experimental results of Kesel and the numerical results of Kim et al., shows good agreement at comparable wing positions [9,11]. Figure 4 shows the corresponding pressure coefficient. The pressure acts vice versa to the velocity and is smaller on the upper side and higher on the lower side. The lift coefficient is the area between the upper and the lower graph of the pressure coefficient. The area is greater considering the pressure coefficient at 5° compared to the area at 0°.

Once the wing cross-sections were rotated about 10° and 15°, the flow around the wing stalled. The results led to periodically oscillating lift and drag coefficients. In Fig. 2a and b the mean values are
Fig. 2. a) Lift and Drag coefficient at the different cross sections (in mm) at angles of attack (AOA) of 0°, 5°, 10° and 15°; At 10° and 15° the mean lift coefficients are chosen; b) Dependency of the L/D-ratio on AOA; At AOA 10° and 15°, the L/D ratio drops.

Fig. 3. Velocity magnitude distribution of the wing cross-section at 0 mm at an AOA of 5°. The velocity magnitude is given in m/s. The highest lift-to-drag ratio occurred at this location.
Fig. 4. Pressure coefficient of the wing cross-section at 0 mm at an AOA of 5°. The x-coordinates are normalized to the chord length of the wing. The graphs of the lower and the upper surface are shown separately.

Fig. 5. Velocity magnitude distribution of the wing cross-section at 5 mm and an AOA of 10° at 0.525 s. The velocity magnitude is given in m/s.

illustrated. The L/D – ratio decreases the higher the AOA. Figure 5 shows the wing cut at 5 mm and an AOA of 10° at 0.525 s. A Kármán vortex street is visible behind the wing. Figure 6 shows the corresponding pressure coefficient distribution as well as the pressure coefficient distribution at the following time steps. The shedding vortices and the time dependency are visible.
5. Conclusions

Due to the steady vortices that occur in the valleys of the corrugated wing profile, recirculation zones are formed. Therefore, the flat wing’s profile resembles the profile of an airfoil. Due to the corrugations, especially on the upper side the pressure decreases more compared to that of an airfoil. This may lead to higher lift coefficients. The corrugations decrease in the direction of the wing tip. The more the wing is corrugated, the higher the drag coefficient. Therefore, the wing might have its lowest drag coefficient at 15 mm. However, the highest drag coefficient results at 10 mm. This could be explained, due to the also high lift coefficient.

After the wing cross-sections were rotated about 10° or 15° the flow started to detach. To resolve the occurring vortices, a very small time-step needed to be selected. The lift coefficient increases the higher the AOA. The drag coefficient increases to a greater extent than the lift coefficient. Therefore, the L/D ratio decreases. Behind the wing a Kármán vortex street is visible. The flow detaches at half of the chord length. Again and again new vortices are formed, which dissolve far behind the wing.

The spikes in the graphs of the pressure coefficients appear when the flow hits the vein structure. The flow stagnates when it hits a vein, therefore the pressure rises. During passing the vein the flow accelerates. The area between the pressure coefficient of the upper and the lower graph at 0 mm is higher at the AOA of 5° compared to the AOA of 0°. This leads to a higher lift coefficient.

The orientation of the wing’s leading edge changes at the wing nodus. From the root to the nodus, the leading edge points upwards, from nodus to tip, it points downwards. This change in the orientation of the leading might influence the lift and drag coefficients. Okamoto et al. states, that the orientation change of the leading edge influences the aerodynamic performance [13]. In this study, especially considering the steady analyses the lift coefficient and the L/D – ratio are higher after the nodus at 5°.

Conflict of interest

None to report.
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