On the oxygen abundance determination in HII regions

High – metallicity regions

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Abstract. This is our second paper devoted to the problem of line intensity – oxygen abundance calibration starting from the idea of McGaugh (1991) that the strong oxygen lines ([OII]λλ 3727, 3729 and [OIII]λλ 4959, 5007) contain the necessary information to determine accurate abundances in HII regions. In the previous study (Pilyugin 2000) the corresponding relations were obtained for the low-metallicity HII regions (12+logO/H ≤ 7.95, the lower branch of the O/H – R23 diagram). The high-metallicity HII regions (12+logO/H ≥ 8.2, the upper branch of the O/H – R23 diagram) are considered in the present study.

A relation of the type O/H=f(P,R23) between oxygen abundance and the value of abundance index R23 introduced by Pagel et al. (1979) and the excitation parameter P (which is defined here as the contribution of the radiation in [OIII]λλ 4959, 5007 lines to the "total" oxygen radiation) has been derived empirically using the available oxygen abundances determined via measurement of a temperature-sensitive line ratio [OIII]4959,5007/[OIII]4363 (T_e - method). By comparing oxygen abundances in high-metallicity HII regions derived with the T_e - method and those derived with the suggested relations (P – method), it was found that the precision of oxygen abundance determination with the P – method is around 0.1 dex (the mean difference for the 38 HII regions considered is ~ 0.08dex) and is comparable to that of the T_e - method.

A relation of the type T_e=f(P,R23) between electron temperature and the values of abundance index R23 and the excitation parameter P was derived empirically using the available electron temperatures determined via measurement of temperature-sensitive line ratios. The maximum value of differences between electron temperatures determined via measurement of temperature-sensitive line ratios and those derived with the suggested relation is around 1000K for HII regions considered here, the mean value of differences for 38 HII regions is ~ 500K, which is the same order of magnitude as the uncertainties of electron temperature determinations in high-metallicity HII regions via measured temperature-sensitive line ratios.

Key words: HII regions; galaxies - galaxies: abundances - galaxies: ISM - galaxies: spiral

1. Introduction

An investigation of chemical properties of galaxies is very important for the development of the theory of structure and evolution of galaxies. Oxygen plays a key role in understanding the (chemical) evolution of galaxies for several reasons. Firstly, good spectrophotometry of HII regions is available now for a large number of galaxies, and the oxygen abundances derived from line intensities are published in many works (Caplan et al. 2000; Deharveng et al. 2000; Esteban et al. 1998, 1999a,b; Garnett et al. 1997; Izotov and Thuan 1998, 1999; Izotov, Thuan, and Lipovetsky 1994, 1997; Kobulnicky and Skillman 1996, 1997, 1998; Kobulnicky et al. 1997; Skillman et al. 1994; van Zee et al. 1997, 1998; among others). Secondly, the origin of oxygen seems to be reliably established in contrast to other elements like carbon or nitrogen. The oxygen abundance can be considered as a tool to investigate the evolution of galaxies. For example, the value of oxygen abundance in a galaxy combined with the value of the gas mass fraction can tell us about the efficiency of mass exchange between a galaxy and its environment (Pilyugin and Ferrini 1998, 2000).

Accurate oxygen abundances can be derived from measurement of temperature-sensitive line ratios, such as [OIII]4959,5007/[OIII]4363. This method will be referred to as the T_e - method. Unfortunately, in oxygen-rich HII regions the temperature-sensitive lines such as [OIII]4363...
are too weak to be detected. For such HII regions, empirical abundance indicators based on more readily observable lines were suggested (Pagel et al. 1979; Alloin et al. 1979). The empirical oxygen abundance indicator $R_{23} = ([\text{OIII}]3727,3729 + [\text{OIII}]4959,5007)/H\beta$, suggested by Pagel et al. (1979), has found widespread acceptance and use for the oxygen abundance determination in HII regions where the temperature-sensitive lines are undetectable. This method will be referred to as the $R_{23}$ - method. Several workers have suggested calibrations of $R_{23}$ in terms of oxygen abundance (Edmunds and Pagel 1984, McCall et al. 1985, Dopita and Evans 1986, Zaritsky et al. 1994, among others).

There are two problems with the oxygen abundances derived by the $R_{23}$ - method. First, oxygen abundances derived with different $R_{23}$ - calibrations can differ by 0.3dex and more. The usually-used $R_{23}$ calibrations (Edmunds and Pagel 1984, McCall et al. 1985, Dopita and Evans 1986, McGaugh 1991) are based on a few then-available oxygen abundance determinations through the $T_e$ - method and HII region models. More oxygen abundance determinations through the $T_e$ - method are available now. None of these $R_{23}$ calibrations can reproduce the available data well enough (Pilyugin 2000, Paper I). Secondly, it has been found (Paper I) that the error in the oxygen abundance derived with the $R_{23}$ - method involves two parts: the first is a random error and the second is a systematic error depending on the excitation parameter. The origin of this systematic error is as follows. In a general case the intensities of oxygen emission lines in spectra of HII regions depend not only on the oxygen abundance but also on the physical conditions (hardness of the ionizing radiation and geometrical factor). Then in the determination of the oxygen abundance from line intensities the physical conditions in the HII region should be taken into account. In the $T_e$ - method this is done via $T_e$. In the $R_{23}$ - method the physical conditions in an HII region are ignored.

In our recent work (Paper I), a new way of oxygen abundance determination in HII regions ($P$ - method) was suggested. A more general relation of the type $O/H = f(P, R_{23})$ is used in the $P$ - method, compared to the relation of the type $O/H = f(R_{23})$ used in the traditional $R_{23}$ - method. It was found in Paper I that the precision of oxygen abundance determinations in low-metallicity ($12 + \log O/H \leq 7.95$) HII regions with the $P$ - method is comparable to that with the $T_e$ - method. The oxygen abundances of high-metallicity ($12 + \log O/H \geq 8.2$) HII regions derived with the $P$ - method are significantly less accurate. This seems to be because a subset of HII regions with high-quality homogeneous determinations of oxygen abundances was used in the construction of the $O/H = f(P, R_{23})$ relation in the case of the low-metallicity HII regions while the relation for high-metallicity HII regions was based on the set of all available HII regions with inhomogeneous determinations of oxygen abundances. Here the oxygen abundances for high-metallicity HII regions with measured line ratios $[OIII]/\lambda 4959,5007/\lambda 3727$ will be recomputed in a uniform manner, and the set of HII regions with homogeneous determinations of oxygen abundances will be used in the construction of the $O/H = f(P, R_{23})$ relation in the case of the high-metallicity HII regions.

The search for an $O/H = f(P, R_{23})$ relation for high-metallicity HII regions is the goal of this study. The preliminary analysis of the relevant observational data and redetermination of oxygen abundance in a uniform way for a large set of HII regions are given in Section 2. The line intensities – $O/H$ calibration is derived in Section 3. The line intensities – $T_e$ calibration is derived in Section 4. A discussion is presented in Section 5. Section 6 contains a brief summary.

## 2. Preliminary analysis of observational data

The strategy for construction of an empirical relation between strong line intensities and oxygen abundance is based on the following propositions 1) the value of $O/H$ in high-metallicity HII region can be expressed as a function of two parameters: the value of $R_{23}$ and the hardness of the ionizing radiation, 2) the excitation index $P$ is a good indicator of the hardness of the ionizing radiation. (Notations similar to those of Paper I will be adopted here: $R_{23} = I_{[OIII]}/(I_{[OIII]} + I_{H\beta}); R_3 = I_{[OII]}/(I_{[OII]} + I_{4363}/I_{H\beta}; R = I_{[OIII]}/I_{H\beta}; R_{23} = R_2 + R_3; R_{23} = \log R_{23},$ and $P = R_3/R_{23}$. The excitation index $P$ used here and indexes $p_2$ and $p_3$ used in Paper I are related through simple expressions: $p_3 = \log P$ and $p_2 = \log(1-P)$. These propositions are immediately evident from the observational data. It has been shown that the value of $R_{23}$ is a robust property of a given HII region (Kennicutt et al. 2000, Oey et al. 2000) in the sense that its value is relatively constant within a given HII region. As an illustration of this fact, Fig. shows the $R_{23} - P$ diagram for the multiple positions in the HII region M17 (Milky Way Galaxy) and in the HII region DEM323 (Large Magellanic Cloud). The M17 data include spectra from Peimbert et al. (1992) and the DEM323 data are from Oey et al. (2000). Both HII regions show a large range of excitation across the regions sampled; at the same time the value of $R_{23}$ is relatively constant. Other examples can be found in Kennicutt et al. (2000) and Oey et al. (2000).

On the other hand, Fig. shows that the value of $R_{23}$ is not a good indicator of metal abundance in HII regions in the sense that there is no one-to-one correspondence between $R_{23}$ and oxygen abundance. Indeed, the HII regions M17 and DEM323 have in fact the same values of $R_{23}$, while their oxygen abundances are rather different: $12 + \log O/H = 8.37$ in DEM323 versus $12 + \log O/H = 8.61$ in M17. Thus the value of $R_{23}$ is governed not only by the oxygen abundance but at least one additional parameter. It is usually accepted that the spectrum of an HII
region (and hence the value of \(R_{23}\)) is specified by three parameters: the abundance of the chemical elements in the gas, the ionization parameter or geometrical factor, and the hardness of the ionizing radiation (Stasinska 1990, McGaugh 1991, among others). The relative constancy of \(R_{23}\) within a given nebula suggests that its value in moderately metal-rich HII region depends very weakly (if at all) on the ionization parameter. This leads us to anticipate that the value of O/H is a function of two parameters: the value of \(R_{23}\) and the hardness of the ionizing radiation. Thus, an important proposition that in moderately high-metallicity HII regions \(R_{23}\) is mainly governed by the oxygen abundance and by the hardness of the ionizing radiation (or effective temperature of the exciting star(s)) is immediately evident from the following observational facts: 1) the value of \(R_{23}\) is relatively constant within a given HII region, 2) there is no one-to-one correspondence between \(R_{23}\) and oxygen abundance.

It has been found that the value of \([\text{OIII}]/\text{H}\beta\) correlates with the effective temperature of the exciting star or the hardness of the ionizing radiation (Kaler 1978, Copetti et al. 1986, among others) and hence the value of \([\text{OIII}]/\text{H}\beta\) can be used as indicator of the hardness of the ionizing radiation. Vilchez and Pagel (1988) have introduced the value

\[
[\text{OII}]\lambda\lambda3727,3729/[\text{OIII}]\lambda\lambda4959,5007 \\
[S\text{II}]\lambda\lambda6716,6731/[S\text{III}]\lambda\lambda9069,9532
\]  

(1)

as a hardness index. \([\text{OIII}]/\text{H}\beta\) is not a very good indicator of the radiation hardness as it is affected by the oxygen abundance as well. The hardness index of Vilchez and Pagel contains the intensity of line \([\text{SIII}]\lambda9532\), which is often unknown. It will be assumed here that the excitation index P can be used as indicator of the hardness of the ionizing radiation. The observational evidence in favor of this proposition is given in Fig. 2, the "calibrating" objects from Kennicutt et al. (2000) show a clearly defined correlation between the excitation index P and effective temperature of the ionizing star.

The spectra of part but not the whole HII region are often observed. It should be noted that the value of P is a good indicator of the hardness of the ionizing radiation only if measured fluxes reflect their relative contributions to the radiation of the whole nebula. This fact, together with the quality of spectra, has been taken into account in compiling the sample of HII regions used for the construction of an empirical relation between strong line intensities and oxygen abundance. If the measurements for individual areas in the same HII region were reported in a paper, the integrated spectra were derived (if authors did not do so). If several spectra of equal quality were available for the same HII region, the spectrum with largest \(\text{H}\beta\) flux was preferred. Our sample includes 38 HII regions with \(12+\log(O/H)>8.2\) for which we have collected the relevant observational data, listed with references in Table I.

It is widely accepted that an accurate oxygen abundance can be derived from measurement of temperature-sensitive line ratios, such as \([\text{OIII}]\lambda4959, 5007/\lambda4363\), i.e. through the \(T_\varepsilon\) – method. In the general case, the oxygen abundances in the same HII region with measured line ratios \([\text{OIII}]\lambda4959, 5007/\lambda4363\) derived in different works can differ for three reasons: atomic data adopted, interpretation of the temperature structure (single characteristic \(T_\varepsilon\), two-zone model for \(T_\varepsilon\), model with small-scale temperature fluctuations) and errors in the line intensity measurements. Therefore the compilation of HII regions with
Table 1. Characteristics for the HII regions in the present sample. The commonly used name of the galaxy is given in column 1, the name of the HII region is reported in column 2. The fluxes $R_2 = I(OIII)\lambda5007/I/H\beta$, $R_3 = I(OIII)\lambda3729/3727/I/H\beta$, $R = I(OII)\lambda3729/I/H\beta$ are listed in columns 3 to 5. The electron concentration $N_e$ is reported in column 6. The source for fluxes and electron concentration is listed in column 7. The electron temperature $T_e$ derived from the measured $R3/R$ line ratio is given in column 8. The electron temperature $T_P$ determined with the suggested $T_e=f(P)$ relation is reported in column 9. The electron temperatures are given in units of 10$^4$K. The oxygen abundance $O/H_T$ derived through the $T_e$ – method with $T_e$ is listed in column 10. The oxygen abundance $O/H_P$ determined via suggested O/H=f(P,R23) relation is given in column 11. The oxygen abundance $O/H_P$, derived through the $T_e$ – method with $T_P$ is listed in column 12. The oxygen abundances are given in units of 12+logO/H.

| galaxy  | HII region | R2   | R3   | R    | $N_e$ | references | $T_e$ | $T_P$ | $O/H_T$ | $O/H_P$ | $O/H_P$ |
|---------|------------|------|------|------|-------|------------|-------|-------|---------|---------|---------|
| MWG     | M16=Sh49   | 1.429| 1.262| 0.0012| 124   | C+D       | 0.66  | 0.63  | 8.58    | 8.68    | 8.67    |
| Sh117   |            | 1.799| 1.506| 0.0018| 17    | C+D       | 0.69  | 0.69  | 8.59    | 8.60    | 8.60    |
| Sh184   |            | 1.900| 1.762| 0.0032| 65    | C+D       | 0.76  | 0.72  | 8.48    | 8.58    | 8.58    |
| Sh206   |            | 1.189| 4.637| 0.0137| 412   | C+D       | 0.85  | 0.80  | 8.47    | 8.57    | 8.58    |
| Sh212   |            | 3.102| 1.741| 0.0068| 126   | C+D       | 0.92  | 0.88  | 8.32    | 8.38    | 8.38    |
| M42=Orion|           | 1.355| 3.912| 0.0101| 357   | C+D       | 0.82  | 0.78  | 8.55    | 8.59    | 8.64    |
| M17     |            | 1.239| 4.720| 0.0100| 691   | P92       | 0.79  | 0.81  | 8.61    | 8.56    | 8.57    |
| S208=RCW5-1|         | 0.912| 9.119| 0.0708| 200   | S83       | 1.13  | 1.06  | 8.27    | 8.34    | 8.36    |
| S311=RCW16-1|       | 2.344| 3.861| 0.0107| 200   | S83       | 0.84  | 0.90  | 8.54    | 8.42    | 8.43    |

List of references:
C+D – Caplan et al. (2000) + Deharveng et al. (2000); D87 – Diaz et al. (1987); G97 – Garnett et al. (1997); I98 – Izotov and Thuan (1998); K96 – Kobulnicky and Skillman (1996); K97 – Kobulnicky and Skillman (1997); O+O – Oey and Shields (2000) + Oey et al. (2000); P80 – Pagel et al. (1980); P92 – Peimbert et al. (1992); S75 – Smith (1975); S83 – Shaver et al. (1983); T89 – Torres-Peimbert et al. (1989); V88 – Vilchez et al. (1988); V97 – van Zee et al. (1997); VZ98 – van Zee et al. (1998); W83 – Webster and Smith (1983).
original oxygen abundance determinations through the $T_e$ – method from different works carried out over more than twenty years is not a set of homogeneous determinations. Accordingly, the available published spectra of HII regions with measured line ratios $[OIII] \lambda\lambda 4959, 5007/\lambda 4363$ (listed in the Table 1) have been reanalysed to produce a homogeneous set. Two-zone models of HII regions with the algorithm for oxygen abundance determination from Pagel et al. (1992) and $T_e([OIII]) - T_e([OII])$ relation from Garnett (1992) were adopted here. The oxygen abundances for HII regions from Table 1 were recomputed in this common way. The obtained electron temperature $T_e$ and corresponding oxygen abundance $O/H$ are reported in columns 8 and 10 of Table 1. The electron temperatures derived by Garnett et al. (1997) were used for NGC2403 HII regions (the $R = I_{[OIII]4363}/I_{H\beta}$ values for NGC2403 HII regions reported in Table 1 are compared the to these electron temperatures of the HII regions and not to the measurements). This set of HII regions with homogeneous determinations will be used in the construction of an empirical relation between line intensities and oxygen abundance.

3. Line intensities – O/H calibration

The $X_{23}$ versus O/H diagram for HII regions from Table 1 is presented in Fig. 3. The open circles are HII regions with $P > 0.5$, the filled circles are HII regions with $P < 0.5$. The line is the $X_{23}$ calibration of Edmunds and Pagel (1984). Inspection of Fig. 3 shows that there is no one-to-one correspondence between $X_{23}$ value and oxygen abundance. For a fixed oxygen abundance the positions of low-excitation HII regions in the $X_{23} - O/H$ diagram are systematically shifted towards lower values of $X_{23}$ compared to the positions of high-excitation HII regions. The calibration of Edmunds and Pagel (1984) corresponds to the positions of high-excitation HII regions. This is not surprising since observational data only for high-excitation HII regions were available when this calibration was suggested. Other previous calibrations (McCall & Rybski and Shields 1985, Dopita & Evans 1986, Zaritsky et al. 1994) are shifted towards still higher oxygen abundances. The fact that for a fixed oxygen abundance the value of $X_{23}$ varies with the excitation parameter $P$ confirms our proposition that the excitation index $P$ can be used in the oxygen abundance determination.

It is convenient to start the search for the relation between oxygen abundance and parameters $P$ and $R_{23}$ from consideration of $R_{3} - P$ (Fig. 4) and $R_{23} - P$ (Fig. 5) diagrams. The positions of HII regions from Table 1 in these diagrams are presented by circles; the filled circles are HII regions with $12 + \log(O/H) < 8.4$, open circles are HII regions with $12 + \log(O/H) > 8.4$. The points are HII regions from Zaritsky et al. (1994) and van Zee et al. (1998). The pluses are low-excitation HII regions in our Galaxy from Deharveng et al. (2000). Fig. 4 shows that HII regions occupy rather narrow band in the $R_3 - P$ diagram. The relation between $R_3$ and $P$ can be given by a polynomial

$$k_0 R_3 = k_1 P + k_2 P^2 + k_3 P^3. \quad (2)$$

The zero-degree term of the polynomial must be equal to zero since $R_3$ and $P$ are equal to zero simultaneously by the definition. Examination of Fig. 4 shows that the positions of the HII regions with $12 + \log(O/H) < 8.4$ are shifted relative to those with $12 + \log(O/H) > 8.4$. In order to take this fact into account the coefficients of Eq. (2) will be taken in the form

$$k_j = a_j + b_j Z, \quad (3)$$

where the notation $Z = 12 + \log(O/H)$ is used for brevity. Taking Eq. (3) into account, Eq. (2) can be rewritten as

$$R_3 = \frac{(a_1 + b_1 Z) P + (a_2 + b_2 Z) P^2 + (a_3 + b_3 Z) P^3}{1 + b_0 Z}. \quad (4)$$

The coefficient $a_0$ has been taken equal to 1 (dividing the numerator and denominator in the right side of Eq. (4) by $a_0$). Eq. (4) can be solved for the value $Z = 12 + \log(O/H)$

$$12 + \log(O/H) P = \frac{R_3 - a_1 P - b_0 R_{23}}{b_1 P + b_2 P^2 + b_3 P^3 - b_0 R_{23}}. \quad (5)$$

Taking into account that $R_3 = P \times R_{23}$, Eq. (5) can be transformed into

$$12 + \log(O/H) P = \frac{R_{23} - a_1 - a_2 P - b_0 R_{23}}{b_1 P + b_2 P + b_3 P^2 - b_0 R_{23}}. \quad (6)$$

The coefficients $b_0, a_1, b_1, a_2, b_2, a_3, b_3$ can be found using the set of HII regions with oxygen abundances derived through the $T_e$ method. In other words, the positions in

Fig. 3. The $X_{23} - O/H$ diagram for the HII regions from our sample. The filled circles are HII regions with $P < 0.5$, open circles are HII regions with $P > 0.5$. The line is the $R_{23}$ calibration after Edmunds and Pagel (1984).
the $R_{23} - P$ (and the $R_3 - P$) diagram can be calibrated in terms of oxygen abundance.

The precision of present-day determinations of the oxygen abundances in high-metallicity HII regions through the $T_e$ - method seems to be around 0.1 dex (Deharveng et al. 2000). It has been found that there are different sets of coefficients $b_0, a_1, b_1, a_2, b_2, a_3, b_3$ (or different variants of the calibration) which give an average value of differences $\Delta O/H_P = \log O/H_P - \log O/H_{T_e}$ less than 0.1 dex for our sample of HII regions. Unfortunately, our sample of HII regions (used in the search for the coefficients) does not contain the HII regions with $P < 0.4$, Fig. 4. Therefore different sets of coefficients based on the narrow range of oxygen abundances in high-excitation HII regions can result in appreciably different abundances in low-excitation ones although all the variants of the calibration that resulted in a correlation between the differences $\Delta O/H_j$ for individual HII regions and $P_j$ or between $\Delta O/H_j$ and $O/H_j$ were rejected and only the variants in which both correlation coefficients are less than 0.1 were considered. Thus, the problem of choice of the variant of the calibration, which results in correct oxygen abundances in the whole range of $P$, appears. This difficulty is usually resolved by adding HII region models to the real HII regions. As was discussed by Stasinska (2000), the existing models of HII regions may be far from reality for a number of reasons. Then we may have to try to overcome this difficulty starting from the observational data only. The following solution to this problem will be adopted. Our sample contains 9 HII regions of the disk of our Galaxy. Fig. 6 shows the oxygen abundance as a function of the galactocentric distance. The linear fit

$$12 + \log(O/H) = 8.90 - 0.041R_G$$  

(7)

is close to the relation obtained by Deharveng et al. (2000). This is not surprising, since the measured fluxes for 6 out of the 9 HII regions were taken from them. The list of HII regions observed by Caplan et al. (2000) and Deharveng et al. (2000) contains a number of low-excitation ones. Following Deharveng et al. (2000) all the heavily reddened HII regions were excluded from consideration. The nine low-excitation HII regions Sh54, Sh131, Sh148, Sh 152, Sh153, Sh156, Sh168, Sh217, and Sh219 for which flux measurements seems to be reliable were selected and used as ”secondary calibrating objects” in the choice of the relation between strong line intensities and oxygen abundance (the positions of these HII regions in the $R_3 - P$ and $R_{23} - P$ diagrams are shown by pluses, Figs. 4, 5). The basic requirement imposed on the calibration by our set of HII regions is complemented by the condition that
the oxygen abundances of selected low-excitation HII regions determined through the T_e – method must result in the same radial oxygen abundance gradient as was determined from HII regions with oxygen abundances derived through the T_e – method. Inspection of Fig.6 shows that positions of HII regions from this "extended" set cover the whole range of P.

The choice of the variant of the calibration which is suitable for the whole range of the parameter P was performed in the following way. A variant of the calibration can be characterized by three parameters: O/H^* = O/H(R_{23}=0.1,P=0), \langle \Delta O/H \rangle = \sqrt{\frac{1}{n} \sum_{j=1}^{n} (\Delta O/H_j)^2}/n, and \langle \delta O/H \rangle = \frac{\sum_{j=1}^{n} \Delta O/H_j}{n} where \Delta O/H_j is equal to \log(O/H)_G^{\prime} - \log(O/H)_{Pj} for our set of HII regions and \Delta O/H_j is equal to \log(O/H)_{Pj} - \log(O/H)_{R_{23}} for the selected low-excitation HII regions where O/H)_{R_{23}} is the abundance estimated through the galactocentric distance (Eq.7). The O/H^* is the minimum value of oxygen abundance for HII regions with R_{23}=0.1, predicted by a given variant of the calibration. The value of \langle \Delta O/H \rangle characterizes the scatter of differences between oxygen abundances determined via a fixed variant of the calibration and adopted oxygen abundances. The value of \langle \delta O/H \rangle characterizes the average systematic difference between oxygen abundances determined via the fixed variant of the calibration and adopted oxygen abundances. For every fixed value of the parameter O/H^* the variant of the calibration which gives the minimum value of \langle \Delta O/H \rangle was obtained. The value of \langle \Delta O/H \rangle_{min} and corresponding value of \langle \delta O/H \rangle as a function of the O/H^* are presented in Fig.7. The local non-smooth variations of the values of \langle \Delta O/H \rangle_{min} and \langle \delta O/H \rangle with O/H^* reflect the fact that all the variants of the calibration which result in a correlation between \Delta O/H_j and P_j or between \Delta O/H_j and O/H_j were rejected and only the variants in which both correlation coefficients are less than 0.1 were considered.

Fig.7 shows that the \langle \Delta O/H \rangle_{min} has a minimum value for O/H^* = 8.91. The corresponding value of \langle \delta O/H \rangle is close to zero (≈ –0.005dex). This variant of the calibration will be referred to as V891. The comparison of this variant of the calibration with observational data is given in Fig.8 (1st column). Two other variants V907 and V915 of the calibration are also presented in Fig.8 (panels in columns 2 and 3). Fig.8 shows that the V891 variant of the calibration well reproduces the observational data of different types. The positions of HII galaxies from our sample in the R_{23} – P and O/H – R_{23} diagrams are well reproduced by the V891 variant of the calibration (panels a1 and b1 in Fig.8). The differences between oxygen abundances determined through the variant V891 of the calibration and through the T_e – method do not exceed 0.1dex (the average difference is about 0.08 dex for 38 HII regions) for the majority of our set (panel c1 in Fig.8). For comparison, the average difference between variant V915 and the T_e – method is in excess of 0.1dex (≈ 0.13dex) (panel c3

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**Fig. 6.** The radial distribution of oxygen abundance in the disk of our Galaxy. The oxygen abundances of HII regions are determined through the T_e – method (recomputed here); the distances are taken from Deharveng et al. (2000).

**Fig. 7.** The \langle \Delta O/H \rangle_{min} and \langle \delta O/H \rangle as a function of the O/H^* = O/H(R_{23}=0.1,P=0). The \langle \Delta O/H \rangle_{min} values are presented by pluses, the \langle \delta O/H \rangle values are presented by the circles. The variants of the calibration (V891, V907, and V915) discussed in the text are shown by filled triangles.
Fig. 8. The comparison of three variants of the calibration (V891 - 1st column panels, V907 - 2nd column panels, V915 - 3rd column panels) with observational data. The panels a1, a2, and a3 show the $R_{23} - P$ diagrams. The filled circles are HII regions with $8.2 < 12 + \log(O/H) < 8.4$, the open triangles are HII regions with $8.4 < 12 + \log(O/H) < 8.62$ from our sample. The pluses are low-excitation HII regions in the disk of our Galaxy from Deharveng et al. (2000). The curves are $R_{23} - P$ relations predicted by a given variant of the calibration for fixed values of $\log(O/H) = 8.2$, $\log(O/H) = 8.4$, $\log(O/H) = 8.6$ (from top to bottom). The panels b1, b2, and b3 show the $X_{23} - O/H$ diagrams. The filled circles are HII regions with $P < 0.5$, the open circles are HII regions with $P > 0.5$ from our sample. The curves are $X_{23} - O/H$ relations predicted by a given variant of the calibration for fixed values of $P=1$, $P=0.5$, and $P=0$ (from top to bottom). The panels c1, c2, and c3 show $\Delta(O/H) = \log(O/H)_P - \log(O/H)_{Te}$ versus $O/H_{Te}$ diagrams for our sample of HII regions. The panels d1, d2, and d3 show the radial distribution of oxygen abundance within the disk of our Galaxy. The circles are HII regions with oxygen abundances derived through the $T_e$ method, the pluses are low-excitation HII regions with oxygen abundances determined through the given variant of the present calibration.
in Fig. 3). Consequently, variant V891 of the calibration results in the radial oxygen abundance gradient which is very close to that derived from HII regions with oxygen abundances determined through the Te – method (panel d1 in Fig. 8). Then, the variant V891 of the calibration is the most credible relation between strong line intensities and oxygen abundance. This variant of the calibration with coefficients \( b_0 = -0.243, a_1 = -54.2, b_1 = 6.07, a_2 = -59.45, b_2 = 6.71, a_3 = -7.31, b_3 = 0.371 \) and corresponding equation

\[
12+\log(O/H)_P = \frac{R_{23} + 54.2 + 59.45P + 7.31P^2}{6.07 + 6.71P + 0.371P^2 + 0.243R_{23}} \quad (8)
\]

can be adopted for oxygen abundance determinations in moderately high-metallicity HII regions with undetectable temperature-sensitive line ratios. The oxygen abundances in HII regions from our sample determined with this calibration are given in Table 1 (column 11).

4. Line intensities – Te calibration

The oxygen abundances \( O/H_P \) derived in the previous section through the excitation parameter \( P \) and abundance index \( R_{23} \) are in agreement with those derived through the classical Te – method. This justifies the use of relations of the type \( O/H = f(P, R_{23}) \) for oxygen abundance determination in high-metallicity HII regions with the undetectable \( [\text{OIII}]4363 \) line. It also provides evidence that the propositions which are at the basis of this relation (i) that the value of abundance index is mainly governed by the oxygen abundance and by the hardness of the ionizing radiation and depends very weakly (if at all) on the ionization parameter, (ii) that the parameter \( P \) can be used as indicator of the hardness of the ionizing radiation) seem to be close to reality. Conversely, if the physical conditions in high-metallicity nebulae are governed mainly by the oxygen abundance and by the hardness of the ionizing radiation and if observational values of \( P \) and \( R_{23} \) reflect these two, then it can be expected that the physical conditions in a nebula can be derived with help of observational values \( P \) and \( R_{23} \). The physical conditions in a nebula are reflected in the electron temperature, so that a relation of the type \( \text{Te} = f(P, R_{23}) \) can be expected.

Fig. 8 shows the positions of HII regions from Table 1 in the \( R_{23} – P \) diagram. The HII regions with \( T_e < 0.8 \) are shown by circles, those with \( 0.8 < T_e < 1.0 \) are shown by pluses and those with \( T_e > 1.0 \) are shown by triangles. Fig. 8 shows that the HII regions with different electron temperatures occupy different bands in the \( R_{23} – P \) diagram. This suggests that the positions in the \( R_{23} – P \) diagram can be calibrated in terms of electron temperatures. Let us again start from the expression of the type

\[
R_3 = \frac{(A_1 + B_1t)P + (A_2 + B_2t)P^2 + (A_3 + B_3t)P^3}{1 + B_0t}, \quad (9)
\]

which can be solved for the value of \( t_P \)

\[
t_P = \frac{R_3 - A_1P - A_2P^2 - A_3P^3}{B_1P + B_2P^2 + B_3P^3 - B_0R_3}, \quad (10)
\]

where \( t_P \) is the electron temperature expressed in units of \( 10^3 \) K. The coefficients \( B_0, A_1, B_1, A_2, B_2, A_3, \) and \( B_3 \) can be found using the sample of HII regions with electron temperature derived via measured temperature-sensitive line ratios \([\text{OIII}]4959,5007/\text{[OIII]}4363 \). It has been found that there are different sets of coefficients \( B_0, A_1, B_1, A_2, B_2, A_3, \) and \( B_3 \) which provide the average value of differences \( \Delta t_P = t_P - T_e \) around 500K. The variant of the calibration (the choice is explained below) with coefficients \( B_0 = 0.583, A_1 = -3.09, B_1 = 9.90, A_2 = -7.05, B_2 = 11.86, A_3 = -2.87, B_3 = 7.05 \) and corresponding equation

\[
t_P = \frac{R_{23} + 3.09 + 7.05P + 2.87P^2}{9.90 + 11.86P + 7.05P^2 - 0.583R_{23}}, \quad (11)
\]

has been adopted for the electron temperature determination in high-metallicity HII regions with undetectable temperature-sensitive line ratios. The electron temperatures \( t_P \) in HII regions from our sample determined from Eq. (11) are given in Table 1 (column 9).

The \( R_{23} – P \) relations predicted by the calibration for fixed values of \( t_P = 0.6, 0.8, 1.0, \) and 1.2 are presented in Fig. 9 by dashed lines. Fig. 10 shows the differences between
electron temperatures \( t_P \) derived through the present calibration and measured electron temperatures \( T_e \). As can be seen in Fig. 10, the largest value of the difference \( \Delta T_e = t_P - T_e \) is around 1000K, the average value about 500K.

Since the electron temperatures \( t_P \) derived through the present calibration are close to the measured electron temperatures \( T_e \), the \( t_P \) can be used for the oxygen abundance determination through the \( T_e \)-method. The oxygen abundances \( O/H_{t_P} \) in HII regions from our sample determined in this way are given in Table 1 (column 12). The differences between oxygen abundances derived with the \( t_P \) and with the measured \( T_e \) are presented in Fig. 11. The variant of the \( t_P=f(P,R_{23}) \) calibration which results in the best agreement between \( O/H_{t_P} \) and \( O/H_{T_e} \) for our set of HII regions and between \( O/H_{t_P} \) and \( O/H_{R_{23}} \) for low-excitation HII regions in the disk of our Galaxy was chosen above.

The differences between \( O/H_{t_P} \) and \( O/H_{T_e} \) for our set of HII regions are presented in Fig. 12. It can be seen in Fig. 12 that \( O/H=f(P,R_{23}) \) and \( t_P=f(P,R_{23}) \) relations lead to consistent abundances.

Thus, the positions in the \( R_{23} - P \) diagram can be calibrated in terms of electron temperatures. The obtained relation \( t_P=f(P,R_{23}) \) between electron temperature and parameters \( P \) and \( R_{23} \) provides an alternative (to the relation \( O/H=f(P,R_{23}) \) obtained in the previous section) method of the oxygen abundance determination in high-metallicity HII regions with the undetectable \([OIII]4363\) line.

5. Discussion

The starting expression (Eq. (2)) for the search for the relation \( T_e = f(P,R_{23}) \) has been chosen from analysis of positions of HII regions in the \( R_{3} - P \) (Fig. 4) and \( R_{23} - P \) (Fig. 5) diagrams. These diagrams are based on directly measured values. The starting Eq. (2) can be rewritten in the form

\[
k_0 R_{23} = k_1 + k_2 P + k_3 P^2,
\]

where coefficients \( k_j \) are dependent on the electron temperature. With our assumption that the parameter \( P \) is an indicator of the hardness of the ionizing radiation, Eq. (12) is the relation between the hardness of the ionizing radiation and electron temperature in the nebula. On the other hand, it is well known that the relation between the hardness of the ionizing radiation (or effective temperature of the exciting star) and electron temperature in the nebula can be derived from the law of energy conservation for
free electrons. According to Sobolev (1967, Eq.23.37) this relation is given by the expression

$$AT_{\text{eff}} = BT_e + C R_{23} + D \left\langle \frac{n_1}{n^+} \right\rangle$$  \hspace{1cm} (13)

where coefficient A depends on $T_{\text{eff}}$ only and B, C and D depend on $T_e$ only. Comparison of Eq.(13) with Eq.(12) shows that with our assumption that the parameter P is an indicator of the hardness of the ionizing radiation (or effective temperature of the exciting star) our starting Eq.(12) is in some sense similar to Eq.(13). Then, the relation $T_e = f(P,R_{23})$ derived here can be considered as some kind analog of the equation of the energy balance of a gaseous nebula.

The validity of the obtained relation $O/H=f(P,R_{23})$ depends on the reliability of the oxygen abundances in HII regions for which this relation has been derived. The two-zone model for $T_e$ is at the basis of the oxygen abundance determination in the present study. Then the $O/H=f(P,R_{23})$ relation cannot be more credible than the two-zone model for $T_e$. If it is ever established that the two-zone model for $T_e$ is a crude approximation of reality and leads to large uncertainty in the oxygen abundances then the oxygen abundances in HII regions should be re-determined within the framework of a more realistic model and $O/H=f(P,R_{23})$ relation should be revised.

Most metal-rich HII regions are objects with small values of $R_{23}$. Although both the O/H$_P$ and the O/H$_{R_{23}}$ are maximum for those HII regions, the O/H$_P$ for them are significantly lower (up to 0.5 dex) as compared to the O/H$_{R_{23}}$. It must be emphasized, however, that the O/H$_P=f(P,R_{23})$ relation is established on the basis of HII regions with $R_{23}$ larger than 2 (panel a1 in Fig 8) and the validity of this relation in the case of HII regions with $R_{23}$ less than 2 may be disputed. Then an additional information should be invoked in order to firmly establish how rich the most oxygen-rich HII regions are.

6. Conclusions

The problem of line intensity – oxygen abundance calibration has been investigated starting from the idea of McGaugh (1991) that the strong oxygen lines ([OII]λ3727, 3729 and [OIII]λλ4959, 5007) contain the necessary information to determine accurate abundances in HII regions. The high-metallicity HII regions (12+logO/H ≥ 8.2, the upper branch of the O/H – $R_{23}$ diagram) are considered in the present study. The low-metallicity HII regions (12+logO/H ≤ 7.95, the lower branch of the O/H – $R_{23}$ diagram) have been considered in the previous study (Pilyugin 2000).

An important proposition that in high-metallicity HII regions the value of the abundance index $R_{23}$ is mainly governed by the oxygen abundance and by the hardness of the ionizing radiation (or by effective temperature of the exciting star(s)) and depends very weakly (if at all) on the ionization parameter, is immediately evident from the following observational facts: 1) the value of $R_{23}$ is relatively constant within a given HII region, 2) there is no one-to-one correspondence between $R_{23}$ and oxygen abundance. Another important fact supported the use of the excitation parameter P as an indicator of the hardness of the ionizing radiation is also evident from the observational data. These propositions are at the basis of the strong oxygen line intensity – oxygen abundance calibration.

A relation of the type $O/H=f(P,R_{23})$ was derived empirically using the available oxygen abundances determined via measurement of temperature-sensitive line ratios ($T_e$ – method). By comparing oxygen abundances in HII regions derived with the $T_e$ – method and those derived with the suggested relations (P – method) it was found that the precision of oxygen abundance determination with the P – method is around 0.1 dex and is comparable to that obtained with the $T_e$ – method.

A relation of the type $T_e=f(P,R_{23})$ was derived empirically using the available electron temperatures determined via measurement of the temperature-sensitive line ratio [OIII]4959,5007/[OII]4363. The maximum value of differences between electron temperatures derived through the $T_e=f(P,R_{23})$ relation and determined via measurement of the temperature-sensitive line ratio is around 1000K for HII regions considered here; the mean difference is ~ 500K, which is of the same order of magnitude as the uncertainties of electron temperature determinations in high-metallicity HII regions via measured temperature-sensitive line ratios themselves.

The relation $T_e$=$f(P,R_{23})$ between electron temperature and parameters P and $R_{23}$ provides an alternative (to the relation O/H=f(P,R_{23})) method of oxygen abundance determination in high-metallicity HII regions with the undetectable [OIII]4363 line. By comparing oxygen abundances in HII regions derived through the $T_e$ – method with measured electron temperatures and those derived through the $T_e$ – method with electron temperature determined with help of $T_e=f(P,R_{23})$ relation, it was found that the precision of oxygen abundance determination with the latter method is around 0.1 dex and is comparable to that obtained with the $T_e$ – method.

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