Photo-induced electric polarizability of Fe$_3$O$_4$ nanoparticles in weak optical fields

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**Abstract**
Using a developed co-precipitation method, we synthesized spherical Fe$_3$O$_4$ nanoparticles with a wide nonlinear absorption band of visible radiation. Optical properties of the synthesized nanoparticles dispersed in an optically transparent copolymer of methyl methacrylate with styrene were studied by optical spectroscopy and z-scan techniques. We found that the electric polarizability of Fe$_3$O$_4$ nanoparticles is altered by low-intensity visible radiation ($I \leq 0.2$ kW/cm$^2$; $\lambda = 442$ and 561 nm) and reaches a value of $10^7$ Å$^3$. The change in polarizability is induced by the intraband phototransition of charge carriers. This optical effect may be employed to improve the drug uptake properties of Fe$_3$O$_4$ nanoparticles.

**Keywords:** Magnetite nanoparticles; Electric polarizability; Low-intensity visible radiation

**PACS:** 33.15.Kr; 78.67.Bf; 42.70.Nq

**Background**
Magnetite (FeO*Fe$_2$O$_3$, or Fe$_3$O$_4$) nanoparticles, and materials based on them, have been successfully used to solve applied problems in biology and magneto-optics. Pronounced superparamagnetic [1-4] and ferromagnetic [5] properties at room temperature enable the use of these nanoparticles in magnetic resonance imaging [6-9] and biosensing [9] as well as in drug delivery and drug uptake applications [8-13]. Because they possess magneto-optical properties [14,15], Fe$_3$O$_4$ nanoparticles have also been used to develop tunable filters [16,17] and optical switches [18,19] that operate under magnetic fields.

In fact, Fe$_3$O$_4$ nanoparticles have been examined for the presence of unique magnetic properties because magnetite is a narrow-gap semiconductor [20-22] and the optical properties of other semiconductor nanoparticles have been thoroughly studied. Currently, there are several experimental and theoretical works dedicated to studying the optical properties of both bulk magnetite [23-26] and its nanoparticles [27-29]. However, some specific optical properties of Fe$_3$O$_4$ nanoparticles (in particular, the effects of electric polarizability on their biological activity, conductivity, ferroelectricity, and electro-optical properties) as well as the nature of these properties remain virtually unexplored.

In this paper, we demonstrate that Fe$_3$O$_4$ nanoparticles exhibiting a wide nonlinear absorption band of visible radiation (1.7:3.7 eV) are able to significantly change their electric polarizability when exposed to low-intensity visible radiation ($I \leq 0.2$ kW/cm$^2$). The observed change in polarizability was induced by the intraband phototransition of nanoparticle charge carriers, and polarizability changes were orders of magnitude greater than those of semiconductor nanoparticles and molecules [30,31].

**Experiments**

**Synthesis of nanoparticles**
There are several techniques for the synthesis of Fe$_3$O$_4$ nanoparticles with an arbitrary shape and size and for their dispersal in different matrices [4,5,11,12,27,29,32-36]. In this study, we synthesized nanoparticles using co-precipitation method [1,2,13-15,37,38], dispersed them in monomeric methyl methacrylate with styrene (MMAS), and polymerized this composition using pre-polymerization method.

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In the first step (Figure 1a), Fe₃O₄ nanoparticles were synthesized by co-precipitation of soluble salts of ferrous and ferric ions with an aqueous ammonia solution: FeSO₄*7H₂O + 2FeCl₃*6H₂O + 8NH₃*H₂O ↔ Fe₃O₄ + 6NH₄Cl + (NH₄)₂SO₄ + 20H₂O. Oleic acid (in a mass ratio of 0.7:1 with the formed Fe₃O₄) was added to a 0.5% solution of iron salts (FeSO₄/FeCl₃ = 1:2.2 molar ratio) in 0.1 M HCl. The aqueous solution of iron salts was heated to 80°C, followed by the addition of concentrated aqueous ammonia (20% excess). The solution was heated and stirred for an hour.

Stabilized nanoparticles were then extracted from the aqueous phase into a nonpolar organic solvent hexane at a ratio of 1:1. The organic layer containing the iron oxide Fe₃O₄ was separated from the aqueous medium. The sample was centrifuged for 15 min (6,000 rpm) to remove larger particles. Excess acid was removed with ethanol.

The size of the nanoparticles was determined by dynamic light scattering method (Zetasizer Nano ZS, Malvern, UK). Measurements were conducted in hexane with a laser wavelength of 532 nm. The average hydrodynamic diameter of the synthesized nanoparticles was 15 nm, as illustrated in Figure 2.

Composite preparation

The second step (Figure 1b) focused on obtaining a solid composite based on Fe₃O₄ nanoparticles and MMAS. The organic solvent containing nanoparticles and monomers (methyl methacrylate with styrene) was subjected to stirring and ultrasonic homogenization. To prevent nanoparticle aggregation during the polymerization process, we used the pre-polymerization method at 75°C because the nanoparticles had different affinities to the monomer and polymer.

Finally, the composite was synthesized in situ by radical polymerization. The polymerization of methyl methacrylate with styrene (in the mass ratio of 20:1) proceeded for over 10 h (in a temperature gradient mode that progressed from 55°C to 110°C) in the presence of benzoyl peroxide (10⁻³ mol/L).

The obtained solid composites had 0.001%, 0.003%, 0.005%, and 0.01% volume concentrations of Fe₃O₄ nanoparticles in MMAS. Importantly, the synthesized Fe₃O₄ nanoparticles generally had a thick layer of acids [36,39] surrounding them to prevent aggregation of the nanoparticle. In our case, the synthesized Fe₃O₄ nanoparticles had a thick layer of acids surrounding them to prevent aggregation of the nanoparticle.

![Figure 1](http://www.nanoscalereslett.com/content/8/1/317)

**Figure 1** The developed co-precipitation method. (a) The synthesis of Fe₃O₄ nanoparticles with a monolayer of oleic acid by the developed co-precipitation method and (b) the composite MMAS + Fe₃O₄ preparation.

![Figure 2](http://www.nanoscalereslett.com/content/8/1/317)

**Figure 2** Nanoparticle size. The average hydrodynamic diameter of the synthesized nanoparticles (15 nm) dispersed in hexane was determined by dynamic light scattering method (Zetasizer Nano ZS, Malvern, UK) at a laser wavelength of 532 nm.
nanoparticles had a monolayer of oleic acid that allowed the nanoparticles to exhibit their specific optical properties.

**UV–vis spectroscopy**

Room-temperature optical absorbance spectra of pure MMAS (Figure 3, black curve) and of the composites were obtained using a Varian Cary 5000I spectrophotometer (Agilent Technologies, Santa Clara, CA, USA) over the wavelength range of 300 to 1,500 nm. These spectra allowed the derivation of the absorbance spectra for Fe₃O₄ nanoparticle arrays (Figure 3, color curves). Figure 3 shows the absorbance values (Abs) and the absorption coefficients ($\alpha = (\text{Abs} \times \ln 10)/l$, where $l = 7.95$ mm is the length of the composite) measured at a maximum radiation intensity of 1 μW/cm².

**z-Scan experiments**

Because they have absorption bands of 380 to 650 nm, Fe₃O₄ nanoparticles should exhibit an optical response upon external radiation with wavelengths in this band [40]. To detect the optical response of the nanoparticles contained in the composite (0.005% nanoparticle volume concentration), we used the standard z-scan technique [41]. This technique enabled the analysis of changes in the absorption coefficient $\Delta \alpha(I)$ and refractive index $\Delta n(I)$ of the composite and pure MMAS, which were induced by weak optical radiation with different intensities 0 to 0.14 kW/cm².

For radiation sources, we used semiconductor lasers of continuous wave (cw) radiation with wavelengths of 442 nm (blue) and 561 nm (yellow) providing maximal intensities of 0.07 and 0.14 kW/cm². Lenses with focal lengths of 75 mm provided the beam waists $\omega_0 = 102$ and 110 μm for blue and yellow radiation (Figure 4b). The length ($L$) of experimental samples of the MMAS and the composite was 2.7 mm (inset in Figure 3).

Because the Rayleigh range $z_0 = \pi \omega_0^2 / \lambda$ exceeded 10 cm, the calculation of $\Delta \alpha$ and $\Delta n$ was performed using the formulae [40,41]:

$$\begin{align*}
\Delta \alpha(I) &= \frac{2\sqrt{2\Delta T(I)}}{L}, \\
\Delta n(I) &= \frac{\lambda \Delta T_{\text{pv}}(I) \times (\alpha + \Delta \alpha(I))}{0.812\pi(1-S)^{0.27}(1-e^{-\alpha+\Delta \alpha(I)L})}.
\end{align*}$$

where $\Delta T(I)$ (Figure 4a) and $\Delta T_{\text{pv}}(I)$ (Figure 5b) were the integral transmitted intensity and the normalized

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**Figure 3** Absorbance spectra for the MMAS and Fe₃O₄ nanoparticle array. The optical absorbance spectra for pure MMAS and Fe₃O₄ nanoparticle arrays with 0.001%, 0.003%, 0.005%, and 0.01% volume concentrations.

**Figure 4** z-Scan results for the MMAS. (a) Curves for z-scans with open (circle) $\tau_0$ and closed (square) $\tau_{\text{pv}}$ apertures at radiation wavelengths of 442 nm (red points, 60 W/cm²) and 561 nm (blue points, 133 W/cm²) for the MMAS sample ($L = 2.7$ mm). (b) Profilometer images for the beam waists $\omega_0$. 
transmittance between the peak and valley at different radiation intensities, respectively; \( \lambda \) and \( \alpha \) were the radiation wavelength and absorption coefficient (Figure 3), respectively, and \( S \) was the fraction of radiation transmitted by the aperture without the sample, which was 0.184.

The experimental curves \( T(I) \) and \( T_{pv}(I) \), which contain information about \( \Delta T \) and \( \Delta T_{pv} \), showed that only the reverse saturable absorption of yellow radiation occurred in pure MMAS (Figure 4a). In contrast, the composite manifested the expected optical response: the shape of the experimental curves \( T(I) \) and \( T_{pv}(I) \) indicated the saturable absorption of visible radiation in the composite and a negative change in its refractive index (Figure 5), and the values of \( \Delta T(I) \) and \( \Delta T_{pv}(I) \) increased linearly with increasing intensities of blue (Figure 5a) and yellow (Figure 5b) radiation.

The approximation of \( T_{pv} \) based on the theoretical curves (solid lines in Figure 5) was performed using the equation [42]:

\[
T = 1 + \frac{2(-p x^2 + 2x - 3p)}{(x^2 + 9)(x^2 + 1)} \Delta \Phi
\]

where the coupling factor \( \rho = \Delta \alpha \times \lambda / 4\pi \times \Delta n \) and the phase shift due to nonlinear refraction \( \Delta \Phi = 2\pi \times \Delta n \times L_{ef} / \lambda \) had the following values: \( \rho = 0.09 \) and \( \Delta \Phi = -0.23 \) and \( -0.5 \) for blue radiation with intensities of 0.019 and 0.054 kW/cm\(^2\) and \( \rho = 0.05 \) and \( \Delta \Phi = -0.7 \) and \( -1.45 \) for yellow radiation with intensities of 0.04 and 0.093 kW/cm\(^2\).

**Discussion**

The saturable absorption of visible radiation with intensities less than 0.14 kW/cm\(^2\) in the composite and the negative change in the refractive index were due to the presence of Fe\(_3\)O\(_4\) nanoparticles since pure MMAS showed only the relatively weak reverse saturable absorption of yellow radiation. Therefore, the experimental data \( \Delta T(I) \) and \( \Delta T_{pv}(I) \) obtained for the composite could be used to calculate the values of \( \Delta \alpha(I) \) and \( \Delta n(I) \) for Fe\(_3\)O\(_4\) nanoparticle arrays (Equation 1), and these values are listed in Figure 6.

Because the observed dependence of \( \Delta n \) on the radiation intensity \( I \) (Figure 6b) for Fe\(_3\)O\(_4\) nanoparticle arrays could be considered a linear function, it can be

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**Figure 5** z-Scan results for the composite. Curves for z-scans with open (circle) \( T(I) \) and closed (square) \( T_{pv}(I) \) apertures at radiation wavelengths of 442 nm (a) (red points, 19 W/cm\(^2\); blue points, 54 W/cm\(^2\); green points, 93 W/cm\(^2\)) and 561 nm (b) (red points, 40 W/cm\(^2\); blue points, 93 W/cm\(^2\)) for the composite sample (\( L = 2.7 \text{ mm} \)) containing Fe\(_3\)O\(_4\) nanoparticle with a 0.005% volume concentration.

**Figure 6** The values of changes in the absorption coefficient, refractive index, and polarizability of Fe\(_3\)O\(_4\) nanoparticles. (a) The dependency of changes in the absorption coefficients \( \Delta \alpha \) of pure MMAS (circle) and Fe\(_3\)O\(_4\) nanoparticle arrays (square and rhombus) on the intensity of radiation with wavelengths of 442 nm and 561 nm. (b) The dependency of changes in the refractive index \( \Delta n \) and polarizability \( \Delta \alpha (\AA^3) \) of Fe\(_3\)O\(_4\) nanoparticle arrays on the intensity of radiation with wavelengths of 442 nm (rhombus) and 561 nm (square); red dashed lines present the contribution of the thermal effect of cw radiation on the change in the refractive index (Equation 3), and blue dashed lines are theoretical approximations based on the approach of free carrier absorption (Equation 4).
assumed that $\Delta n$ was caused by the thermal effect of the radiation. We estimated the contribution of this effect to the changes of the composite refractive index using the equation [43]:

$$\Delta n_{\text{therm}} = \frac{\Delta E \times dn}{c_{\text{hc}} \rho_d}, \quad (3)$$

where $c_{\text{hc}}$ was the MMAS heat capacity (0.7 J/g·K), $\rho_d$ was the MMAS density (1.3 g/cm$^3$), $dn/dT$ was the MMAS thermo-optic coefficient ($-10^{-5}$ K$^{-1}$), and $\Delta E$ was the energy absorbed by the composite per unit volume per second. The thermal effect of cw low-intensity radiation on the change in the refractive index (red dashed lines in Figure 6b) was relatively small (not more than 20% for blue radiation and 8% for yellow radiation).

Generally, the possibility of a nonthermal optical response of the composite due to external optical radiation is associated with the polarization of Fe$_3$O$_4$ nanoparticles in the external field $E$. Nanoparticle polarization occurs at the spatial separation of positive and negative charges, i.e., at the electron transition to higher allowed energy states (quantum number $l \neq 0$). These transitions should be accompanied by the absorption of external radiation. In our case, we observed the absorption of radiation with wavelengths of 380 to 650 nm (Figure 3). This absorption band consisted of three maxima (380, 480, and 650 nm), indicating the broadened quantum-size states for the electrons in Fe$_3$O$_4$ nanoparticles. Because the bandgap of magnetite is rather small (approximately 0.2 eV) [20-22], the conduction band of Fe$_3$O$_4$ nanoparticles may be employed to significantly improve the drug uptake properties of Fe$_3$O$_4$ nanoparticles.

The amplitude of the nanoparticle polarization is determined by $|E|$ of the external field and the nanoparticle susceptibility ($\chi$) or polarizability ($\alpha$) measured in cubic angstrom. In turn, the change in the refractive index induced by the radiation is associated with the change in nanoparticle polarizability $\Delta \alpha$ (Å$^3$) by classical relations [48]. Therefore, we could calculate the values of $\Delta \alpha$ (Å$^3$) for Fe$_3$O$_4$ nanoparticle using the experimental values of $\Delta n(I)$ and the following equations (SI):

$$\begin{align*}
\epsilon &= n^2(I) - k^2(I) = 1 + \chi \\
\Delta \chi &= \Delta \alpha \left(\frac{\lambda^3}{\pi}\right) \times 10^{30} \cdot N [m^{-3}] 
\end{align*} \quad (5)$$

where $\epsilon$ was the real part of the dielectric constant, the composite refractive index $n(I) = n_0 + \Delta n(I)$, and $n_0$ was the refractive index of pure MMAS (approximately 1.5). The extinction coefficient $k = \alpha \lambda / 4\pi r$ was significantly less than $n(I)$ and could be ignored; $\chi$ was the nanoparticle susceptibility, and $N$ was the nanoparticle concentration (approximately $2.3 \times 10^{19}$ m$^{-3}$). Therefore, the values of $\Delta \alpha$ (Å$^3$) for Fe$_3$O$_4$ nanoparticle were calculated using the formula $\Delta \alpha (\lambda^3) = 2\pi \times \Delta n(I) \times 10^{30} / N$ and are presented in Figure 6b.

The obtained values for the changes in nanoparticle polarizability are orders of magnitude greater than those for semiconductor nanoparticles and molecules [30,31] in extremely weak optical fields. In addition, the average nanoparticle volume was approximately 2.2 $\times$ 10$^{-6}$ Å$^3$, and the maximum value of $\Delta \alpha$ (Å$^3$) was $9 \times 10^6$ Å$^3$. Thus, we can conclude that the nanoparticle polarization should be formed by several optical intraband transitions of nanoparticle electrons in weak optical fields.

**Conclusions**

We used the developed co-precipitation method to synthesize spherical Fe$_3$O$_4$ nanoparticles covered with a monolayer of oleic acid that possessed a wide nonlinear absorption band of visible radiation 1.7 to 3.7 eV. The synthesized nanoparticles were dispersed in the optically transparent copolymer methyl methacrylate with styrene, and their optical properties were studied by optical spectroscopy and z-scan techniques. We report that the electric polarizability of Fe$_3$O$_4$ nanoparticles changes due to the effect of low-intensity visible radiation ($I \leq 0.2$ kW/cm$^2$; $\lambda = 442$ and 561 nm) and reaches a relatively high value of $10^7$ Å$^3$. The change in polarizability is induced by the intraband photon transition of charge carriers and can be controlled by the intensity of the visible radiation used. This optical effect observed in magnetic nanoparticles may be employed to significantly improve the drug uptake properties of Fe$_3$O$_4$ nanoparticles.

**Abbreviations**

Abs: Absorbance; Cw: Continuous wave; MMAS: Methyl methacrylate with styrene.
Competing interests
The authors declare that they have no competing interests.

Authors’ contributions
VM designed and performed the optical experiments (z-scan and spectroscopy), participated in the analysis and interpretation of data, and prepared the draft and final version of the manuscript. AN, VC, and VS designed and performed the chemical experiments, achieved that nanoparticle was covered with a monolayer of oleic acid, prepared the sections ‘Synthesis of nanoparticle’ and ‘Composite preparation’. YK and VD conceived of the study, participated in the analysis and interpretation of data, helped to draft the final version of the manuscript. All the authors read and approved the final manuscript.

Acknowledgements
The work was supported by the Programs of Presidium of Russian Academy of Science (12-1-OFN-05, 12-I-P24-05, 12-II-UO-02-002) and by the Program of UB RAS (12-S-2-1004).

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Received: 23 May 2013 Accepted: 1 July 2013

Published: 9 July 2013

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