Infrared spectroscopy around 4 μm of Seyfert 2 galaxies: Obscured broad line regions and coronal lines

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Abstract. The state of the matter that is obscuring a small circumnuclear region in active galactic nuclei can be probed by observations of its broad emission lines. Infrared lines are particularly useful since they penetrate significant columns of obscuring matter, the properties of which can be constrained by comparing infrared and X-ray obscuration. We report on new 4 μm spectroscopy with ISAAC at the ESO VLT of a sample of 12 Seyfert 2 galaxies, probing for broad components to the Brackett α 4.05 μm hydrogen recombination line. Broad components are observed in 3 to 4 objects. All objects with a broad component exhibit relatively low X-ray obscuring columns, and our results are consistent with a Galactic ratio of 4 μm obscuration to the BLR and X-ray column. In combination with observations of a non-Galactic ratio of visual obscuration of BLRs and X-ray obscuring column in Seyferts, and interpreted in a unified AGN scheme, this result can be reconciled with two interpretations. Either the properties of dust near the AGN are modified towards larger grains, for example through coagulation, in a way that significantly flattens the optical/IR extinction curve, or the ratio of dust obscuration to X-ray column varies for different viewing angles with respect to the axis of symmetry of the putative torus. Our spectra also provide a survey of emission in the [Si IX] 3.94 μm coronal line, finding variation by an order of magnitude in its ratio to Brα. The first extragalactic detection of the [Ca VII] 4.09 μm and [Ca V] 4.16 μm coronal lines is reported in the spectrum of the Circinus galaxy.

Key words. galaxies: active – galaxies: Seyfert – galaxies: ISM

1. Introduction

Unified scenarios have been highly successful in explaining several aspects of the AGN phenomenon, by assuming that different manifestations of the AGN phenomenon correspond to similar objects viewed from different directions. The detection in polarized light of broad emission lines in Seyfert 2 galaxies (Antonucci & Miller 1985; Antonucci 1993) has been central to the development of these scenarios. Subsequently, spectropolarimetry has become the prime tool for detecting hidden broad lines in larger samples (e.g., Miller & Goodrich 1990; Tran et al. 1992; Young et al. 1996; Heisler et al. 1997; Moran et al. 2000; Lumsden et al. 2001; Tran 2001). X-ray spectroscopy has been the second key observation, finding Seyfert 2s on average much more highly obscured than Seyfert 1s and quantitatively establishing the absorbing column densities in neutral and “warm”, i.e. highly ionized material (e.g., Turner et al. 1997; Bassani et al. 1999). Still, relatively little is known about the actual distribution and physical state of the material obscuring the central engine of Seyfert 2 galaxies from our view. Is it in the form of a compact parsec scale torus (Krolik & Begelman 1986)? Or does obscuration on scales of tens or hundreds of parsecs play a significant role (e.g. Maiolino & Rieke 1995; Malkan et al. 1998)? Mass arguments suggest that at least the Compton-thick absorbers are on scales of tens of parsecs or less (Risaliti et al. 1999), but lower column components are more difficult to constrain. Do outflows contribute to forming the obscuration (e.g. Königl & Kartje 1994; Elvis 2000)? How does the state of the obscuring matter differ from the interstellar medium in a normal galaxy, given the extreme conditions close

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to a powerful AGN? High obscuring columns can be explained by different scenarios, and the warm dust emission seen in the mid-infrared cannot uniquely distinguish between compact and more extended configurations either (e.g., Pier & Krolik 1992; Efstathiou et al. 1995; Granato et al. 1997).

One way to address some of these issues is to study the obscuration of the central engine in wavelength ranges that are partially transparent at the column densities of interest, and compare the results. Of particular value are X-rays and the infrared range, but care has to be taken to compare obscuration towards similar regions. Narrow Line Region (NLR) emission and mid-infrared dust emission probe regions that are much larger than the central X-ray source. Therefore, their obscuration may involve different foreground material, and cannot be compared directly to X-ray results. In contrast, reverberation mapping results show the Broad Line Region (BLR) to be well below a parsec in size (Netzer 1990), allowing a meaningful comparison of BLR and X-ray obscuration, by material that might be found in both parsec-scale and larger regions. Our goal is to constrain infrared obscuration, and thus the state of the obscuring matter by the detection or nondetection of infrared BLRs in Seyfert 2s of various X-ray obscuring columns. Recently, Maiolino et al. (2001a) have compared visual and X-ray obscuration in Seyferts, finding large deviations from standard Galactic values. In many objects, the ratio of reddening $E_{B-V}$ and X-ray column $N_H$ is low by about an order of magnitude compared to the Galactic value. These results provide additional motivation to explore the relation between infrared and X-ray obscuration, and compare it with that in the visual.

Practical considerations drive the choice of the optimal transition for Broad Line Region searches in the infrared. Standard interstellar extinction laws have a minimum in the 3–8 $\mu$m range, then rise through the silicate features and drop again steeply beyond 30 $\mu$m. This tends to argue in favour of the longest wavelength infrared observations. However, the flux of the strongest ($\alpha$) recombination lines drops approximately with the second power of wavelength, while dust continuum increases making their detection increasingly difficult. On the basis of such reasoning and of ISO spectroscopy of Brackett $\beta$, Brackett $\alpha$, and Pfund $\alpha$ in NGC 1068, Lutz et al. (2000a) concluded that Brackett $\alpha$ with its fairly high line to continuum ratio and low obscuration is the most promising line for infrared BLR searches. Using sensitive instruments such as ISAAC on 8 m telescopes like the VLT, it is now possible to perform such observations and take a next step beyond earlier infrared searches for BLRs, which either focused on lines that have a good line to continuum but still relatively high obscuration ($Pa\beta 1.28 \mu$m) or are suffering less extinction but are difficult to measure because of a low line-to-continuum ratio (Bry 2.17 $\mu$m). Searches in these lines detected several broad components, but cannot probe beyond equivalent visual obscurations of 10–20 mag (e.g., Rix et al. 1990; Blanco et al. 1990; Goodrich et al. 1994; Ruiz et al. 1994; Veilleux et al. 1997; Gilli et al. 2000).

This paper is organized as follows. Section 2 discusses the observations and data analysis, Sect. 3 the result of a coronal line survey obtained from our data, Sect. 4 presents the results of the Brackett $\alpha$ spectroscopy including line decompositions. We discuss the results and compare to X-ray and optical work in Sect. 5 and conclude in Sect. 6.

### 2. Observations and data analysis

To identify suitable bright Seyfert 2 galaxies we have used the Bassani et al. (1999) sample with well-known X-ray obscuring columns. We have selected the objects that are brightest in extinction-corrected [O III] $\lambda 5007 (>10^{-12}$ erg s$^{-1}$ cm$^{-2}$), observable from Paranal (declination $<20$ deg), and are at $z < 0.015$ to keep Brackett $\alpha$ in the observable part of the atmospheric $L$ band. Of the 23 objects in this parent sample, 12 were observed following right ascension and weather constraints during a run in March 9–11, 2001. Both the parent sample and the observed targets cover a wide range of X-ray obscuring columns. NGC 1068 was part of the parent sample and is included in our discussion, using the existing ISO data of Lutz et al. (2000a, 2000b) observed with a large $14'' \times 20''$ aperture. Due to the prominence of the central $L$ band peak in NGC 1068 (Alonso-Herrero et al. 1998), the continuum of these data is still dominated by the central compact peak and can be analysed together with the smaller aperture ISAAC sample. For part of our sample, spectropolarimetric observations are available in the literature, directly confirming the intrinsic presence of a Broad Line Region. Table 1 lists this information, together with some basic source properties and total line fluxes measured from the spectra.

We used the ISAAC-LW medium resolution spectroscopy mode to cover a range of 3.93 to 4.17 $\mu$m at spectral resolving power $\approx 2500$. This range was chosen because the rapid drop of atmospheric transmission makes longer wavelengths useless, and to ensure coverage of the [Si IX] 3.94 $\mu$m coronal line in all spectra. While we present below results for this and other coronal lines obtained from our spectra, the primary motivation for including this line was to provide a reference, observed with the same instrument setting, for the line widths in the Narrow/Coronal Line Region. These can be used in cases where there is ambiguity between a true BLR and relatively broad NLR components. The bright and compact nuclei were acquired in the $K$ band and centered in the $1''$ slit which was oriented north-south. Observations were done in the chopping and nodding along the slit scheme suited for thermal infrared observations. Integration times (excluding overheads) varied between 37 and 62 min per source. Our strategy was to integrate to good signal-to-noise in the narrow component of Br$\gamma$. Then, nondetection of a broad component is meaningful, since in typical Seyfert 1s the broad line flux is several times brighter than the narrow line flux. In the sample of Stirpe (1990), for example, the median flux ratio of broad and narrow H$\beta$ is 20 and the lowest ratio 7, for those 14 objects that do not exceed a luminosity of $M_{V} = -22$ mag.

Data reduction followed standard procedures in Eclipse and IRAF. The wavelength calibration is based on arc spectra and on a vacuum scale. For flux calibration and correction for atmospheric absorptions, we observed both early type and G type stars. For our particular case of spectra near Brackett $\alpha$ with sometimes good continuum S/N, we found early type standards
little suited because of difficulties to correct for their intrinsic hydrogen and helium absorption lines, and in some cases emission lines. More satisfactory correction of atmospheric absorption was achieved using early G star spectra which had first been corrected for their significant intrinsic spectral structure using the high resolution solar spectrum provided at the ESO ISAAC web pages. This solar spectrum was convolved to the appropriate resolution and slightly shifted and scaled, to give optimum cancellation of G star features when dividing a G star and solar spectrum. Since we are interested mostly in the inner region showing dust and BLR rather than NLR emission we used optimum extraction of the spectra to get the best...
signal-to-noise. Effects on the NLR fluxes which in principle occur were verified to be small by inspection of the 2-D spectra and comparison of optimum extracted with directly extracted spectra. Only Circinus showed evidence for faint extended Br\(\alpha\) emission peaking \(\approx 10-15''\) from the nucleus, which we did not include in the extracted spectrum. Figures 1 and 2 show the spectra, first as observed and then sorted by the X-ray obscuring column and scaled in a way facilitating comparison. Total fluxes of the coronal lines and of Br\(\alpha\) (from direct integration of the line profiles) are listed in Table 1. Due to the narrow slit and optimum extraction scheme, they correspond to a small \(\approx 1''\) aperture.

3. A survey of coronal lines

The [Si IX] 3.94 \(\mu m\) line was first observed in an extragalactic object by Oliva et al. (1994) in the spectrum of the Circinus galaxy. It is one of the highest excitation coronal lines observed
in the spectra of Seyfert galaxies, with a lower ionisation potential of 303 eV. It is thus an interesting diagnostic of the coronal line region of AGN, the density of which can be constrained to be less than about $10^6$ cm$^{-3}$ using the density-sensitive ratio [Si IX] 2.58/3.94 µm (e.g. Lutz et al. 2000b). Our good signal-to-noise 4 µm spectra of 12 Seyferts significantly widen the database of observations of this line. Figure 3 shows the distribution of the ratio of [Si IX] 3.94 µm and the narrow (NLR) component of Br$\alpha$, where we have again added NGC 1068 (Lutz et al. 2000b) to our sources. While the [Si IX] line is detected in 12 of 13 objects, the observed excitations differ strongly, with the ratio of [Si IX] and Br$\alpha$ varying by about an order of magnitude, similar to variations in the strength of optical coronal lines in Seyferts (e.g., Penston et al. 1984; Erkens et al. 1997).

Factors that might bias the observed ratio with respect to the intrinsic NLR one need consideration. Firstly, Br$\alpha$ may be contaminated by starburst emission. This effect is likely to be small in our small aperture. This can be tested on the basis of the line profiles. On the high excitation end, for Circinus both Br$\alpha$ and [Si IX] lines are consistent with the literature radial velocity of Table 1 (from NED). As expected for our resolving power which is relatively low for detailed NLR profile studies, both lines are consistent with the literature radial velocity in most objects. The significant blueshift of high excitation lines in NGC 1068 has been discussed in detail by Lutz et al. (2000b). In MCG-5-23-16, both lines are redshifted by ~80 km s$^{-1}$, possibly indicating an inaccuracy of the optical-based NED redshift. NGC 3081 appears to be a case like NGC 1068, with blueshift and asymmetry of the coronal line persisting in the infrared. An interesting case deserving further study is NGC 4507, where a redshift of the coronal line with respect to Br$\alpha$ is suggested.

We conclude that these effects do not dominate the [Si IX]/Br$\alpha$ spread of about an order of magnitude among the objects of our sample. We also computed simple CLOUDY photoionization models for the “table power law” continuum implemented in CLOUDY and varying ionization parameter, showing that already modest variations of $\Delta$$\log(U)$ ~ 0.3 near $\log(U) = -2$ can change [Si IX]/Br$\alpha$ by an order of magnitude. The observed [Si IX] flux is hence a good qualitative indicator for the presence of an AGN, but a relatively poor quantitative indicator of the AGN luminosity, due to the considerable variations of observed excitation.

Two more coronal lines are detected for the first time to our knowledge in the spectrum of an extragalactic source: [Ca VII] 4.09 µm and [Ca V] 4.16 µm in the spectrum of the Circinus galaxy, at wavelengths in agreement with the ones inferred by Feuchtgruber et al. (2001) from observation of the planetary nebulae NGC 6302 and NGC 7027. Like [Si IX] and Br$\alpha$, they show a distinct blue asymmetry, the blue wing extending out to ~800 km s$^{-1}$. A hint of the [Ca VII] line is seen in other high excitation spectra (NGC 4388, MCG-5-23-16) but at low significance. The rest wavelength of [Ca V] is covered by our spectra for Circinus only. Detection of these lines (lower ionisation potentials 109 and 67 eV) confirms again the unique suitability of the Circinus galaxy for coronal line studies due to its brightness, high excitation, and narrow line widths.

Observations of infrared fine-structure and coronal lines can also help elucidating the role of dust in forming the line asymmetries and shifts that are frequently observed in optical coronal lines (e.g., Erkens et al. 1997). For the prototypical Seyferts NGC 1068 and NGC 4151, these issues are discussed on the basis of ISO data by Lutz et al. (2000b) and Sturm et al. (1999). They observe some of the differences between optical and infrared lines that are expected for a moderate amount of dust in the NLR, but also find remaining asymmetries and shifts in the infrared that must either reflect intrinsic asymmetry of the NLR, or an optically extremely thick absorber that suppresses even infrared lines. We list in Table 2 the velocity offsets of narrow Br$\alpha$ and [Si IX] with respect to the radial velocity of Table 1 (from NED). As expected for our resolving power which is relatively low for detailed NLR profile studies, both lines are consistent with the literature radial velocity in most objects. The significant blueshift of high excitation lines in NGC 1068 has been discussed in detail by Lutz et al. (2000b). In MCG-5-23-16, both lines are redshifted by ~80 km s$^{-1}$, possibly indicating an inaccuracy of the optical-based NED redshift. NGC 3081 appears to be a case like NGC 1068, with blueshift and asymmetry of the coronal line persisting in the infrared. An interesting case deserving further study is NGC 4507, where a redshift of the coronal line with respect to Br$\alpha$ is suggested.

![Fig. 3. Histogram for the ratio of [Si IX] 3.94 µm and the narrow component of Br$\alpha$ in 13 Seyfert galaxies.](image-url)
Table 1. Source properties and integrated line fluxes. Heliocentric redshifts are from NED, extinction corrected [O III] fluxes from Bassani et al. (1999). X-ray obscuring columns \(N_{\text{H}}\) are from Bassani et al. (1999) except for Circinus (Matt et al. 1999).

| Source     | \(cz\) km s\(^{-1}\) | \([\text{O III}]\) 10\(^{-14}\) W m\(^{-2}\) | \(N_{\text{H}}\) 10\(^{20}\) cm\(^{-2}\) | BLR in polarized light? | \([\text{Si IX}]\) | \([\text{Br} \gamma]\) 10\(^{-20}\) W m\(^{-2}\) | \([\text{Ca VII}]\) | \([\text{Ca V}]\) |
|------------|----------------|-----------------|----------------|----------------|----------------|-----------------|----------------|----------------|
| NGC 2992   | 2311           | 0.680           | 69             | yes (Lumsden et al. 2002) | 873             | 6650            |                |                |
| MCG-5-23-16 | 2482           | 0.409           | 162            | yes (Lumsden et al. 2002) | 2150            | 5010            |                |                |
| NGC 2110   | 2335           | 0.321           | 289            | <500             | 2110            |                |                |                |
| NGC 5506   | 1853           | 0.600           | 340            | ? (Tran 2001, Young et al. 1996) | 1000:11700      |                |                |                |
| NGC 4507   | 3258           | 0.138           | 2920           | yes (Moran et al. 2000) | 2380            | 3650            |                |                |
| NGC 388    | 2524           | 0.374           | 1200           | yes (Young et al. 1996) | 2530            | 3220            | ?              |                |
| NGC 4941   | 1108           | 0.355           | 4500           | no (Moran et al. 2000) | 455             | 804             |                |                |
| NGC 3081   | 2385           | 0.215           | 6600           | yes (Moran et al. 2000) | 1500            | 1120            |                |                |
| Mkn 1210   | 4046           | 0.482           | >10\(^4\)      | yes (Tran et al. 1992) | 1230            | 2680            |                |                |
| NGC 4968   | 2957           | 1.116           | >10\(^4\)      |                | 667             | 1060            |                |                |
| Circinus   | 449            | 6.970           | 410\(^4\)      | yes (Oliva et al. 1998) | 37700           | 13900           | 4040           | 4610           |
| NGC 5643   | 1199           | 0.694           | >10\(^5\)      | no (Moran et al. 2000) | 870             | 2920            |                |                |
| NGC 1068   | 1137           | 15.86           | >10\(^5\)      | yes (Antonucci & Miller 85) | 54000           | 69000           |                |                |

Table 2. Kinematic data and \([\text{Br} \gamma]\) line decompositions from Gaussian fits. For objects without BLR detection, the \(FWHM\) of a fit with a single Gaussian is given. Total fluxes from the fits may deviate from the values in Table 1 which are based on a direct integration of the profile. \(\Delta v\) describes the shift of the Gaussian fitted (narrow) line centroid with respect to the galaxy’s radial velocity as listed in Table 1.

| Source     | \(FWHM\) \((b)\) km s\(^{-1}\) | \(FWHM\) \((n)\) km s\(^{-1}\) | \(\Delta \nu\) \(\Delta \nu\) km s\(^{-1}\) | Flux \((b)\) 10\(^{-20}\) W m\(^{-2}\) | Flux \((n)\) 10\(^{-20}\) W m\(^{-2}\) | \([\text{Br} \gamma]\) \(FWHM\) \([\text{Si IX}]\) cm\(^{-2}\) | \(\Delta \nu\) \([\text{Si IX}]\) km s\(^{-1}\) |
|------------|--------------------------------|--------------------------------|--------------------------------|-----------------|-----------------|-----------------|----------------|
| NGC 2992   | 1850                          | 230                           | -23                           | 5370            | 1300            | 240             | -14            |
| MCG-5-23-16| 1450                          | 230                           | +84                           | 3930            | 1110            | 190             | +77            |
| NGC 5506   | 618                           | -20                           | <3100                        | 7150            | 5260            | 640:150         |                |
| NGC 4507   | 1217 fixed: 460              | +39                           | <4000                        | 520             | +74             |                |                |
| NGC 4388   | 540                           | -13                            | <1600                        | 190             | -5              |                |                |
| NGC 4941   | 260                           | -22                            | <3100                        | 130             | +16             |                |                |
| NGC 3081   | 260                           | +016                           | <2300                        | 550             | -89             |                |                |
| Mkn 1210   | 570                           | -1                             | <3900                        | 670             | +16             |                |                |
| NGC 4968   | 270                           | -26                            | <2300                        | 400             | -31             |                |                |
| Circinus   | 240                           | -26                            | <20000                      | 170             | +6              |                |                |
| NGC 5643   | 250                           | -35                            | <3100                        | 290             | +4              |                |                |
| NGC 1068   | 640                           | -27                            | <80000                      | 590             | -187            |                |                |

4. Results of the Brackett \(\alpha\) spectroscopy

In clear cases like NGC 2992, the broad component of Brackett \(\alpha\) is easily discriminated from the narrow component. At intermediate line widths, the situation is sometimes ambiguous, however. It is not obvious whether a line originates in a true BLR of dense clouds close to the AGN, or whether it represents an unusually wide NLR profile that is seen similarly in the forbidden and coronal lines (see the example of NGC 1068; e.g. Lutz et al. 2000a, 2000b). We hence applied the definition of a broad line as “a component of Brackett \(\alpha\) with \(FWHM\) around 1000 km s\(^{-1}\) or more that is not seen in the forbidden or coronal lines”, and used the \([\text{Si IX}]\) line as reference. Table 2 summarizes our decompositions from Gaussian fits. In cases where we do not detect a broad component, we quote an approximate upper limit for its flux determined under the assumption of a representative value of 3000 km s\(^{-1}\) \(FWHM\) for the BLR (e.g. Osterbrock 1977). One should note, however, that the spectral coverage of the data makes detection of very broad \(FWHM\) 10000 km s\(^{-1}\) lines difficult even if the flux were larger.

A number of objects deserve an individual justification of our decision to identify a line as originating in a BLR or not, also to indicate remaining ambiguities. The Brackett \(\alpha\) line of Mkn 1210, for example, can be plausibly decomposed into a broad and a narrow component. We decided against a BLR interpretation, however, because the \([\text{Si IX}]\) line shows wide wings as well, the two profiles being indistinguishable within the S/N limitations. This argues against the interpretation of Veilleux et al. (1997) who considered Mkn 1210 a BLR detection (but see caveats in their appendix). Their [Fe II] and Paschen \(\beta\) profiles are very similar to our \([\text{Si IX}]\) and Brackett \(\alpha\) profiles, while we do not find evidence for the \(FWHM\) 3000 km s\(^{-1}\) component possibly seen in their Brackett \(\gamma\) data.

Another complex case is NGC 5506 where a broad component to Paschen \(\beta\) has been reported (Blanco et al. 1990; Rix et al. 1990). Goodrich et al. (1994) and Veilleux et al. (1997) confirm the larger width of this infrared line compared to optical lines but, on the basis of a similar [Fe II] profile, ascribe this effect to a moderately broad but obscured component of the NLR. Our profile of Brackett \(\alpha\) is relatively broad but not obviously the sum of a narrow and a broad line. Comparison to
[Si IX] does not help in this case because of its weak and uncertain detection in NGC 5506, making the measured FWHM highly uncertain. We hence compared the Brackett α profile with ISO spectroscopy of the [O IV] 26 μm forbidden line in this object (Sturm et al. 2002) which is able to penetrate large obscuring columns. We fixed the NLR FWHM for the ISAAC spectrum to the value of 460 km s\(^{-1}\) which is based on the FWHM measured from a Gaussian fit to [O IV], corrected for the resolution difference of ISO-SWS and ISAAC. Adopting this NLR line width we obtain a residual BLR component with FWHM similar to the one of Blanco et al. (1990) but significantly lower than the value reported by Rix et al. (1990), most likely reflecting line profile decomposition uncertainties of the various datasets. We also checked for differences between the fine-structure line FWHM (Sturm et al. 2002) and the optical FWHM of Veilleux (1991a, 1991b) that would be expected if extinction within the NLR dominates the profiles. No significant differences were found when convolving the optical data to the lower resolution and consistently measuring the FWHM by gaussfits. On the basis of an O I fluorescence line and of J band [Fe II] transitions, Nagar et al. (2002) argue that NGC 5506 is an obscured Narrow Line Seyfert 1, consistent with our decomposition which ascribes most of Brα to a fairly narrow BLR.

The most uncertain case is NGC 2110 with its moderately wide Brackett α line. The nondetection of [Si IX] and the absence of ISO spectroscopy makes a direct comparison to the forbidden lines impossible. Evidence on possible broad components to Paschen β is mixed (Rix et al. 1990; Veilleux et al. 1997). We have listed this case as a narrow line, on the basis of the similarity of the Brackett α line width with that of [Fe II] (Veilleux et al. 1997). We will include this object as uncertain when discussing the sample statistics below.

Figure 4 summarizes the detections and limits on broad components of Brackett α in our sample objects. Relatively few broad line regions are found – only 3 to 4 out of 13 (considering the uncertain case of NGC 2110). For the nondetections, limits for broad Brα components are of the order 1 to 2 times the narrow Brα flux. Assuming intrinsic presence of a BLR, as confirmed by spectropolarimetry for most of the objects, and an intrinsic broad to narrow line ratio 20 equal to the median of the Stirpe (1990) M V < −22 Seyfert 1s, these limits imply an obscuration of broad Brackett α of ≳3 mag. For a Galactic extinction curve, this corresponds to an equivalent visual extinction of more than 50 mag. For A\(_{1.05}/A_V = 0.035\) (Draine 1989) the equivalent visual extinction would be 86 mag, for A\(_{1.05}/A_V = 0.051\) (Lutz 1999) 59 mag. Assuming a canonical conversion factor from visual extinction to X-ray obscuring column of N\(_H/A_V = 1.79 \times 10^{21}\) cm\(^{-2}\) (e.g. Predehl & Schmitt 1995), broad components to Brackett α should remain visible up to X-ray obscuring columns of about 10\(^{23}\) cm\(^{-2}\). The fact that all our BLR detections are below this limit is thus consistent with a Galactic value of the ratio of 4 μm obscuration to X-ray column.

At this point, a comparison to the factors determining the detectability of Seyfert 2 BLRs in polarized light is in place. The dominant factor for BLR detectability in polarized light is AGN luminosity, probably both through the corresponding variation of host galaxy dilution and through a varying size of the scattering region (e.g. Alexander 2001; Lumsden et al. 2001; Lumsden & Alexander 2001). The BLR detectability in polarized light one modestly depends on the X-ray column (Lumsden et al. 2001), there are BLR detections up to the highest X-ray columns. Since most of our objects do have detections of the BLR in polarized light (Table 1), and were selected to be bright in [O III], we believe that the effect of AGN luminosity and host dilution on the BLR detectability in our spectroscopy is less important than the effect of obscuration.

5. Discussion

Our observations have resulted in the detection of broad line regions in about one quarter of the Seyfert 2 sample studied. This fraction is similar to the one reported at shorter near-infrared wavelengths by Veilleux et al. (1997). Small numbers and the difficulties discussed above to ascribe moderately broad profiles for some of the objects to a BLR or the NLR limit a detailed intercomparison of these fractions. It is clear, however, that the detection rate at the wavelength of 4 μm which samples columns up to 10\(^{23}\) cm\(^{-2}\) (for a Galactic extinction curve) is not much higher than in the shorter wavelength studies which penetrate 3 to 5 times smaller columns. This cannot reflect intrinsic absence of BLRs since they are detected through spectropolarimetry in the majority of the objects with 4 μm BLR nondetections. The majority of Seyfert 2s thus still has considerable BLR obscuration at 4 μm. This finding is still consistent with most popular models of dust around AGN, though being close to constraining the ones that predict the most modest obstructions (e.g., Granato et al. 1997). It is already exceeding the modest obscurations (A\(_V\) up to a few tens) predicted for most polar angles (θ < 80 deg) in the disk wind model of Königl & Kartje (1994).

The high 4 μm obscuration and the consistency with a Galactic extinction curve raise an apparent conflict with the abnormally low ratios of visual reddening E\(_{B−V}\) and X-ray column in luminous AGN, as summarized by Maiolino et al. (2001a). Comparing reddening E\(_{B−V}\) towards the BLR as measured from optical spectra with X-ray N\(_H\), they find this ratio to be 3 to 100 times lower than Galactic for most of the objects in their sample. Based on other evidence, they suggest the ratio of visual extinction A\(_V\) and N\(_H\) to be lower than Galactic as well, for several classes of AGN. Simply lowering the dust content in the absorber and releasing the metals into the gas does not solve this discrepancy of about an order of magnitude. In addition to other effects discussed by Maiolino et al. (2001b), a low dust-to-gas ratio should result in the detection of broad Brackett α lines at X-ray columns approaching 10\(^{24}\) cm\(^{-2}\), but none are observed in our sample.

Geometry effects may play a role since the samples are different. Maiolino et al. (2001a) mainly observed Seyfert 1s, Quasars and intermediate Seyferts, since optical detection of a BLR was pre-requisite for their analysis. Here, we analyze intermediate and type 2 objects, that is ones viewed more “edge-on” in the picture of a central torus. If, as not implausible, dust were preferentially modified or destroyed along the opening of the torus, then different lines of sight probe different dust
The seemingly discrepant optical and 4 $\mu$m results are, however, also consistent with the preferred interpretation of Maiolino et al. (2001b) for the results of Maiolino et al. (2001a). If dust in the dense circumnuclear regions of AGN is dominated by large grains, as proposed by Laor & Draine (1993) and by Maiolino et al. (2001b), the net effect on the extinction curve will be mainly a reduction at short wavelengths, but much less change in the infrared. This is illustrated in the top panel of Fig. 5, where the Galactic “standard” extinction curve is compared with an extinction curve due to a grain distribution biased in favor of large sizes. This effect would explain both the mismatch between optical extinction and gaseous column measured in the X-rays (discussed in Maiolino et al. 2001a) and the agreement between IR extinction and X-ray absorption found in this paper.

The “large grains” curve has a turnover at about 2–3 $\mu$m (at variance with the “Galactic” curve) which is nicely probed by the three infrared hydrogen lines Pa$\beta$, Br$\gamma$ and Br$\alpha$. In case of (screen) absorption, the deviation of the line ratios from the intrinsic case B depends significantly on the extinction curve, this is shown in the lower panel of Fig. 5. For two of the sources showing broad Br$\alpha$, namely NGC 2992 and MCG-5-23-16, there are measurements in the literature for the broad components of Pa$\beta$ and Br$\gamma$ (Gilli et al. 2000; Veilleux et al. 1997), which allow to locate these objects in the lower panel of Fig. 5. The location of NGC 2992 is inconsistent with the Galactic extinction curve, but fully consistent with the “large grains” scenario. For MCG-5-23-16 the location is not even consistent with the Galactic extinction, because of the very low Br$\alpha$/Pa$\beta$ ratio. This might be due to variability of the BLR, since Br$\alpha$ (this paper) and Pa$\beta$ (Veilleux et al. 1997) were not measured simultaneously but with a time lag of 9 years. This point of view is strengthened by the data of

More specifically, the distribution of grain sizes has been modelled with $n \propto a^{-\beta}$, $\beta = 2.5$, $a_{\text{min}} = 0.005 \mu$m, $a_{\text{max}} = 1 \mu$m, among those suggested by Laor & Draine (1993) and by Maiolino et al. (2001b), at variance with the “standard” distribution for the Galactic dust, where $\beta = 3.5$, $\alpha_{\text{min}} = 0.005 \mu$m, $a_{\text{max}} = 0.25 \mu$m. See also Maiolino et al. (2001b).
Blanco et al. (1990) who, another 3 years earlier, measured a 4 times lower broad Paβ flux. The caveat of variability might also apply to NGC 2992, though here the measurements were much closer in time (2 years) and the data from Gilli et al. (2000) have higher accuracy. Additional, simultaneous measurements of the broad components of the infrared hydrogen lines are required to unambiguously test the “large grains” scenario.

We note that a population of large grains in the dense circumnuclear regions of AGN is not an ad hoc requirement, but in line with the flattened extinction curves in Galactic dense molecular clouds (e.g. Cardelli et al. 1989) which are ascribed to grain coagulation. However, grain growth may yield the formation of complex fluffy aggregates (e.g., Dominik & Tielens 1997) rather than the simply larger but spherical grains, which are assumed in the simplified model of Fig. 5. Finally, a dust distribution biased for large grains may also come from the preferential destruction of small grains by sublimation or sputtering in the circumnuclear region of AGN. Depending on the mechanism producing a dust distribution that is biased toward large grains, the absolute extinction $A_4 \mu m/N_H$ may differ. In case of destruction of small grains it will be below but similar to the Galactic value, while for coagulation at a fixed dust mass it will rise above the Galactic value. Figure 4 is consistent with modest deviations in either direction.

The current data do not allow to discriminate between the “geometrical” and the “large grains” scenario. Both are plausible within the unified scenario for AGN. A better understanding, at least for the intermediate type Seyferts with moderate $A_V$ could be gained through simultaneous observations of the main optical to 4 $\mu m$ recombination lines, in order to directly trace the flattening of the extinction curve expected in the large grain scenario. Line ratios formed by the broad components of Brackett α, Brackett γ, and Paschen β (Fig. 5) are better suited than a similar diagram invoking the Balmer decrement, because of the reduced susceptibility to departures from case B (e.g. Netzer 1990 and references therein). While Fig. 5 proposes a direct test for the presence of large grains, future quasi-simultaneous and accurately calibrated data will be needed for a conclusive result.

6. Conclusions

We have presented new 4 $\mu m$ spectroscopy of a sample of 12 Seyfert 2 galaxies that are well-studied in the X-ray, and combine these data with previous spectroscopy of NGC 1068. The observations are designed to probe for the presence of optically obscured Broad Line Regions. The main results are: (i) Broad components to Brackett α are detected in 3 to 4 of these 13 galaxies. (ii) The detections and limits are consistent with a Galactic ratio of infrared and X-ray obscuring columns. This result can be reconciled with the low ratios of optical to X-ray obscuring columns observed for several AGN if either the obscuring dust consists of large grains leading to a modified extinction curve, or if variation in dustiness or dust properties exists between directions probing right through the putative torus and directions closer to its opening.

(iii) A survey of the coronal [Si IX] 3.94 $\mu m$ line shows considerable variation in its ratio to Brackett α. (iv) Two coronal lines of [Ca VII] and [Ca V] are detected for the first time in an extragalactic object, the Circinus galaxy.
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