Evolutionary population synthesis: the effect of binary systems

J.M. Mas-Hesse, M. Cerviño
LAEFF-INTA, P.O. Box 50727, E-28080 Madrid, Spain.

Abstract.

We present in this contribution our set of multiwavelength synthesis models including the evolution of single and binary stars. The main results we have obtained can be summarized as follows: (a) massive close-binary systems will start to experience mass transfer episodes after the first 4 Myr of the starburst evolution; (b) as a result of these mass transfer processes, stars of relatively low initial mass can lose completely their envelope and become a WR. In this way, the formation of WR stars is extended over longer than 15 Myr, and does not stop at 6 Myr as predicted by models including only single stars; (c) WR stars can thus be coeval with red supergiants, which peak at around 10 Myr for solar metallicities; (d) the accretion of mass will originate relatively massive stars at ages for which they should have already disappeared; these stars, together with the WR stars formed in rather evolved clusters, increase the production of ionizing photons, so that the Hβ equivalent width will not drop as rapidly as predicted by models considering only individual stars; and (e) the mass transfer to compact companions will produce an additional source of high-energy radiation in the form of high-mass X-ray binaries, not predicted either by standard synthesis models.

1. Introduction

In the last years convincing observational evidences have been collected about the presence of starburst regions in or around active Seyfert 2 nuclei (Heckman et al. 1997; González-Delgado et al. 1998, and references therein). It has been found that most of the UV light from these objects originates in (circum-)nuclear star formation sites; the possible connection between the nuclear activity and the properties of these starbursts is nevertheless still a matter of debate. According to the unified schemes of Seyfert galaxies, the active nucleus should be hidden by an opaque torus in the case of Seyferts 2, explaining so why the UV light collected is indeed dominated by these young, massive stars. The low degree of contamination by the active source in the UV range allows therefore to get detailed information about the properties of the star formation processes. Moreover, extrapolating to the radio – X-ray ranges the emission associated to the starbursts, it should be possible to disentangle the fractional contribution of both sources (the starburst and the active nucleus) at different energy ranges.
Multiwavelength evolutionary synthesis models normalized to the observed UV emission would be the ideal tool to perform this analysis.

With these ideas in mind we started some years ago a program to extend our evolutionary synthesis models to high energy ranges (soft and hard X-rays, gamma-rays). Preliminary results were presented in Cerviño et al. (1996). It was clear that to properly reproduce the high energy emission, the effects of mechanical energy release (by stellar winds and supernova explosions) on heating the diffuse interstellar gas had to be included, as well as the evolution of binary systems. The process of mass transfer between close interacting companions can affect the evolution of the individual stars, and can also trigger high energy emission when one of the stars has already evolved to a compact object. Binaries have also been included in starburst computations by Vanbeveren et al. (1997) and Van Bever & and Vanbeveren (1998). Schaerer & Vacca (1998) performed also a simple approximation to the effects of binary systems on the evolution of young starbursts. In a recent review (Vanbeveren et al. 1998) the detailed description of the evolution of massive close binary systems has been presented. This has also been summarized by Vanbeveren et al. in several contributions in this volume.

We present in this contribution our set of stellar population models including the evolution of single and binary stars. We will stress the effects of binary evolution on the stellar population structure. In Sect. 2 we present our set of evolutionary synthesis models and in Sect. 3 we discuss the predictions on the stellar population. In an accompanying contribution in this volume (Cerviño & Mas-Hesse, “Hard X-ray-to-radio energy distributions in starburst galaxies”), the multiwavelength spectral energy distribution and the effects on the H$\beta$ equivalent width and HeII emission line are discussed. The complete set of models will be published elsewhere.

2. Evolutionary synthesis models for single and binary stars

The basic ideas about our evolutionary synthesis models for single stars have been already discussed in Mas-Hesse & Kunth (1991, hereafter MHK) and in Cerviño & Mas-Hesse (1994, hereafter CMH). Basically, our models are based on different Initial Mass Function (IMF) slopes ($\alpha = 1, 2.35 - $Salpeter– and 3), two extreme star formation regimes (Instantaneous –IB– and Extended Bursts –EB–) and five sets of stellar evolutionary tracks with different metallicities taken from the Geneva group (Schaller et al. 1992 and references therein). In this section we discuss the framework we have used to include the evolution of binary systems, which is based on the prescriptions of Vanbeveren (1991) and Vanbeveren et al. (1998), somewhat modified and simplified.

We assume that the stars evolve like single, non-rotating stars following their theoretical tracks until mass transfer episodes or supernova explosions occur. The procedure for track interpolation in mass has been discussed in detail in CMH. We want to remark that we compute the evolution of each individual mass along the HR diagram, and obtain at each time step the total number of stars and their luminosity for each spectral type and luminosity class bin. In case of the stars forming a binary system, we recompute continuously their mass ratio and orbital separation using their instantaneous masses, as defined
by the stellar tracks. We have assumed circular orbits along the evolution of
the systems, although we have also considered the possible formation of very
eccentric systems after a supernova explosion.

2.1. Initial distributions

The parameters determining the evolution of binary systems are mainly the
mass of each component, their mass ratio \( q = M_2/M_1 \) where \( M_1 \) corresponds
to the more massive star) and the orbital semi-axis, \( a \) or period \( P \), together with
the fraction of stars formed in such systems. The initial distributions of these
parameters will define the evolution of the young clusters and their observational
properties.

| Input parameter        | Law                                      | Method        |
|------------------------|------------------------------------------|---------------|
| IMF                    | \( dN/dm \propto m^{-\alpha} \) (\( \alpha = 1, 2.35, 3 \)) | Montecarlo    |
|                        | Mass limits: 2-120 \( M_\odot \)         |               |
| Star Formation Regime  | Instantaneous, extended                  |               |
| Binary frequency       | 50\% (10, 30, 70, 90\%)                 |               |
| Mass ratio             | \( q = M_2/M_1 \)                        | Montecarlo    |
| Orbital separation     | \( dN/da \propto a^{-1} \)              | Montecarlo    |
|                        | \( a_{\text{min}} < 6(M_1/M_\odot)^{1/2} \), \( a_{\text{max}} = 5000 \) |               |

Table 1. Initial input parameters distributions used in the models.

We have assumed by default that around 50\% of the stars are formed in
binary systems, supported by the frequency of WR stars in binaries obtained
by Vanbeveren & Conti (1980) for the solar neighbourhood. Nevertheless, we
have also explored the effect of assuming different binary frequencies (10, 30, 70
and 90\%). In general, even assuming low frequencies around 10\%, the effects
on the evolution of the stellar population are already remarkable. We want to
stress that this binary frequency \( f \) refers to the total number of binary systems,
and not just to those experiencing mass transfer. Binary interactions will affect
indeed to only less than 5\% of all systems (Maeder & Meynet 1994).

The Initial Mass Function is generated by a Montecarlo technique, as explained
in Arnault (1990), which provides an initial mass value for each star
formed in the cluster. It is forced to follow a power-law with 3 fixed values of the
slope. The mass ratio distribution has been obtained from the same Montecarlo
run, by randomly associating two subsequent mass values, only constrained to
comply with the defined binary frequency. Finally, the orbital separation distribution
has been also generated by a Montecarlo routine, forced to a \( dN/da \propto a^{-1} \)
function. In order to have statistical significance, we have used a distribution
of 5\times10^5 stars (i.e. 1.25\times10^5 binary systems). As explained in CMH, smaller
distributions of stars can lead to significant stochastic variations in the stellar
population from one Montecarlo run to another. While this feature reproduces
what we observe in small clusters, for the purposes of this work we have assumed
only large numbers of stars with well defined initial input distributions.
The initial distributions we have assumed are summarized in Table 1.
2.2. Roche Lobe Over-Flow phase (RLOF)

Mass transfer episodes during Roche lobe overflow are the most critical process when analyzing the evolution of binary systems. We test continuously whether RLOF episodes occur by comparing the Roche lobe size with the stellar radii at each time step. Vanbeveren et al. (1997) show that systems in which the primary has an initial mass larger than 40 M$_\odot$ with periods such that RLOF should start during or after the LBV phase, can avoid the mass transfer episode due to the loss of mass via stellar winds, so that no interaction will occur. We have therefore neglected possible mass transfer episodes in stars initially more massive than 40 M$_\odot$ or when the primary has reached the WR phase when RLOF should occur. This restriction has important consequences, in the sense that the properties of a cluster containing binary systems will not differ significantly from a single-stars-only cluster during the first Myrs of evolution, until the most massive stars in the clusters explode as supernovae and might originate binary systems with a compact companion. When RLOF occur, the evolution of the donor and gainer remnants will depend mostly on their initial mass, their evolutive status, and the amount of mass transferred:

- **Donor star.** The donor star will generally loss its outer layers during mass transfer episodes. Depending on its structure and evolutive status we can identify 3 main cases:
  
  i. **The Donor is an H burning star (Case A).** In this situation mass transfer will last until H is exhausted in the nucleus of the donor star, so that the transfer episode will be extended over the remaining main sequence lifetime of the star. As a result, the donor lifetime will be longer that that of a single star with the same initial mass. As a net effect of Case A episodes the mass ratio of the system will be reversed, and the individual stars will continue their evolution according to their new mass, considering that
around 50% of the mass has been assumed to get lost during the transfer episode. We show in Fig. [1a] the evolution of a system experiencing Case A mass transfer.

ii. Donor between H and He burning phase (Case B). In this case, the star looses mass until the atmospheric H abundance is around 0.2 (de Greve & de Loore 1992). The remnant mass will be, according to Vanbeveren et al. (1998):

\[ m_{\text{post RLOF}}^{d} = 0.093 \times (M_{\text{pre RLOF}}^{d})^{1.44} \]  

As a first approximation we have assumed this law for all metallicities considered, but it may change with metallicity as discussed by Vanbeveren et al. (1998). If the remnant mass of the donor star is larger than 5 M\(_{\odot}\), it becomes a WR-like star, according to Vanbeveren et al. (1997). Bare He burning stars formed after RLOF and with less than 5 M\(_{\odot}\) have been counted also as Wolf-Rayet stars, although their expected weaker, but non-zero, stellar winds might not be able to form a thick mantle originating the spectral features associated with WRs. The value of this mass limit will be refined in the future, when more precise atmosphere models for this kind of stars become available. In the meantime we want to stress that the WR population our models are predicting at evolved ages (longer that 11 Myr) does not necessarily show the same spectral features than more massive WRs.

In order to follow the evolution of the star after mass transfer, we have assumed that: the evolutive status remains the same (no time correction); the surface abundances are equal to the core ones; the instantaneous mass is recalculated at each time step taking into account the remnant mass and the mass loss laws for the different WR phases; the WNL phase will last until all the H is removed from the envelope, assuming a mass loss rate of \(4 \times 10^{-5} M_{\odot} \text{ yr}^{-1}\) and the WC phase will follow the mass-luminosity and mass-radius relations from Schaerer & Maeder (1992) for WR stars and Langer (1989) for He stars (mass after RLOF lower than 5 M\(_{\odot}\)). In Fig [1b] we show the evolution of a system experiencing a Case B RLOF episode.

iii. Donor in He burning or later phases (Case C). In this case the effects of shell-burning regions in the structure of the star is not clear, so that the evolution after RLOF is quite uncertain. Anyway, since the donor remnant will be already very evolved, its remaining lifetime will be very short and the number of these systems will be a minority (Vanbeveren et al. 1997). We have therefore not considered this case of mass transfer episodes in our models.

- Gainer star. We have assumed that the gainer star will be able to accrete the mass lost by the donor if it is still in the Main Sequence, with a convective envelope. In such cases we have considered that 50% of the mass and angular momentum of the system is lost anyway by the system in the mass transfer process (Meurs & van den Heuvel 1989). The gainer will become a more massive
star, which will evolve as a younger one. Combining the evolutive status of the star before the RLOF episode and the amount of mass accreted, we compute a time correction and make the star to evolve following the track of a younger, more massive star. The net effect of these processes will be the “rejuvenation” of the cluster (Van Bever & Vanbeveren 1998), which will host massive stars at ages where all of them should have already exploded as SN, if only single stars were present.

If the gainer is in more evolved phases when RLOF occurs, we have assumed that it will not be able to accrete mass, and that all the mass transferred is lost by the system. In these cases common envelope evolution and spiral-in processes might become important.

2.3. SN explosions and formation of High Mass X-Ray Binaries

After a supernova explosion, the binary system will survive only if the mass of the non-exploding star is two times larger than the mass of the exploding one. We test this conditions when each of the massive stars in the synthetic cluster exhaust its nuclear fuel, and compute then the number of systems that will remain bounded or unbounded after the SN explosion. If the exploding star had at the moment of the explosion a mass higher than 8 M_⊙, we assumed that it becomes a black hole with more than 4 M_⊙. If the mass is lower than 8, but higher than 4 M_⊙, it will become a neutron star with around 3 M_⊙. Otherwise, it will become a White Dwarf with 1.4 M_⊙. After the SN explosion the system might become very eccentric, depending on its mass ratio and orbital separation before the explosion. We have assumed that this eccentric post-SN systems might become a Be/X star (highly eccentric systems with a Main Sequence star plus a compact companion). As the remnant star evolves, it will become a Supergiant. Then we assume that the orbit will be circularized, giving rise to a “permanent” High Mass X-Ray Binary (HMXRB).
Figure 3. (a) Predicted WR population for two metallicities. (b) WR over WR + RSG ratio. (c) WC over WR ratio. Predictions for continuous star formation rate (CSFR) have also been included.

3. Model predictions

We will summarize shortly in this section the most relevant predictions of our evolutionary models affected by the presence of binary systems.

- Wolf-Rayet stars population. One of the most relevant effects of the mass transfers processes in close binaries is the generation of WR stars at ages where models based on single stars only do not predict the presence of any WR at all. We show in Fig. 2 the predicted number of WRs, as well as their types and ratios over O stars, as a function of the binary frequency. It can be seen that while the generation of WR stars from the evolution of single stars remains essentially the same, after around 5 Myr a relatively large number of WR stars is originated in close binary systems. We want to stress that massive stars in detached binary systems will evolve essentially as single stars, and that this new population of WR stars is related only to the binary systems experiencing interactions.

In Fig. 3 we show the dependence of the WR population properties on metallicity. It is interesting to see that the formation of WRs by the binary channel is also metallicity dependent, as a consequence of the different evolution of single stars at lower metallicities. An important result of our models is the prediction of the simultaneous presence of WR and Red Supergiant (RSG) stars in evolved (older than 7 Myr) clusters formed at metallicities above around Z=0.008.

Finally, it is also important to see that our models predict that the WR stars formed by the binary channel will be preferentially of the WC type. As explained above, WR stars formed in Case B will lose their envelope and will show a surface abundance similar of that in the nucleus, i.e., rich in Carbon.

- OB stars and \( W(H\beta) \). The accretion of mass by transfer episodes will originate continuously a population of relatively massive stars at ages where they should have already disappeared from the young clusters (assuming nearly instantaneous bursts). These stars, together with the hot WR stars formed, will maintain the ionizing power of the clusters at a significantly higher level than just single stars. As a result, the strength of the emission lines will decrease...
with time much slower. With a binary frequency of 50%, \(W(H\beta)\), frequently used as an age estimator, will be larger by more than an order of magnitude at 15 Myr than the value predicted for single stars clusters.

- **Systems with compact objects.** The fast evolution of the massive primary originates the formation of close systems with a compact (black hole, neutron star) component after the first 4–5 Myr of evolution. The accretion of mass onto the surface of the compact companion will give rise to high-energy phenomena. As an example, we predict around 10 HMXRB to be present in a 15 Myr cluster having formed around \(10^7\) M\(_\odot\) of stars.

- **Spectral Energy Distribution (SED).** The SED will be drastically affected by the presence of binary systems in the cluster. The formation of HMXRB and related objects provides a source of high-energy emission which completely dominates the hard X-ray range (above around 2 keV) after the first 4–5 Myr. More details are given in the accompanying contribution by Cerviño & Mas-Hesse (this volume).

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