Research Article

Boiling Heat Transfer and Critical Heat Flux Enhancement Using Electrophoretic Deposition of SiO₂ Nanofluid

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The electrophoretic deposition (EPD) technique was used to create a uniform SiO₂ thin film coating on boiling plates, 4 mm in width and 9 mm in length. Significant enhancement in critical heat flux (CHF), for the hydrophilic surfaces generated by this anodic EPD method, has been observed. In order to increase the coating strength, the plates were sintered at various temperatures. To find the thickness and uniformity of the coatings, the SEM images were captured. The captured images showed that the coating thickness uniformly increased up to 90 nm for 0.5% nanofluid percentage by the EPD method. The results show that the hydrophilic and super-hydrophilic surfaces have different boiling heat transfer (BHT) coefficients and CHF behaviors. Also, the results showed an increase of 160% in the CHF value by sintering compared to a bare surface. However, because of the setup simplicity, the shape independency, the particle-coating uniformity, and thickness controllability, the EPD technique can be an appropriate option for modification of the surface and coating on the nuclear fuel cladding.

1. Introduction

One of the efficient heat transfer modes in the nuclear power plants is the nucleate boiling. In this heat transfer regime, the high heat flux can be transferred with smaller wall superheat than other heat transfer regimes. The nucleate boiling regime in the heat transfer systems is limited by the onset of nucleate boiling (ONB) and critical heat flux (CHF). The CHF phenomenon is a transition from the nucleate boiling with a high heat transfer to the film boiling by which heat is removed from the surface by vapor, whereby the heat transfer coefficient is rapidly decreased and the surface temperature is sharply increased with system failure. Therefore, the physical phenomenon of the CHF and its enhancement is an important part in the development of the nuclear thermal hydraulics.

Recent studies of boiling heat transfer have shown that the nanofluids can significantly improve the CHF. These nanofluids are the nanoparticles such as metal oxides, allotropes of carbon, or metals dispersed in base fluids [1]. CHF and boiling heat transfer coefficient have been investigated in the nanofluids boiling heat transfer studies [2–4], whereby these phenomena have changed with the dilute nanofluids. Further studies have shown that these changes in the nanofluid boiling phenomenon and CHF are not due to the intrinsic nature of the nanofluids but are due to the deposition of the nanoparticles on the heating surfaces and change of surface texture and morphology [5–8].

Enhancement of the CHF and boiling heat transfer by surface modification is not a new concept. The major focus on the surface modification to augment CHF and boiling heat transfer has been on the microscale modifications [9–11]. These microstructure modifications have some disadvantages, for example, they reduce thermal stress and thermal resistance. To reduce these microstructure disadvantages of the surfaces as well as to impose an effective control over the surface roughness and porosity, the nanoscale modifications have been recommended.

In this study, to create a nanocoating on the surfaces, the EPD has been used. The EPD technique is utilized in the novel processing of painting, coating, and advanced ceramic
materials [12]. Recently, because of its cost-effective quality and simple apparatus as well as its high versatility to be used with different materials, this contrivance has gained ever-increasing attention in industry and academic investigations. In this technique, the charged nanoparticles dispersed in a base fluid are exposed to DC electric field, and then they are attracted and deposited onto one of the electrodes. Particles with negative charge deposited on the positive electrode (anode) are termed as the anodic EPD, and for positive charge particles, there is cathodic deposition.

In this study, the deposition of SiO$_2$ nanoparticles with anodic EPD is investigated. The deposition was scrutinized on stainless steel in various nanoparticle depositions, voltage ranges, and deposition times. The thickness of the deposition and the composition of the layer are analyzed by a scanning electron microscope (SEM) and electron energy loss spectroscopy (ELS).

2. Experimental Procedure

2.1. Preparation of Nanofluid. SiO$_2$ nanoparticles with deionized water were used to produce nanofluids. Charlot et al. [13] measured a constant high $\xi$ potential ($-$30 ± 5 mV) with a pH of 10 for SiO$_2$ nanofluid. Therefore, to achieve this PH, the alkali medium was obtained using a NaOH aqueous solution. A stable and monodisperse suspension required for the Kline and McClintock [14] uncertainty analysis method is used as follows:

$$q'' = h_{eff} (T_W - T_{sat}).$$  \hspace{1cm} (2)

The pool temperature was adjusted by the heater. The heat flux of the test sample was increased by step increasing of the power supply voltage. This condition continued until the test sample failed; in other words, the CHF was observed at this point. The test sample heat flux was calculated by

$$q'' = \frac{IV}{A}. \hspace{1cm} (3)$$

The boiling heat flux and test sample temperature were recorded to plot the boiling curve and calculate the heat transfer coefficient in every step of the power loading.

To investigate the uncertainty of the CHF measurement, the Kline and McClintock [14] uncertainty analysis method is used as follows:

$$\frac{w_f}{q'} = \left(\frac{w_f}{T}\right)^2 + \left(\frac{w_c}{V}\right)^2 + \left(\frac{w_A}{A}\right)^2. \hspace{1cm} (4)$$

The uncertainty for voltage and current was typically 1% and for cross section was 1.1%. The uncertainty in the thermal flux is 1.8%. Measurement uncertainty for wall superheat was 15%.

3. Results and Discussion

3.1. Contact Angle Measurements. The contact angle is usually measured, where a solid surface contacts with a liquid droplet, to determine the wettability of that surface. The relative molecular interaction strength between vapor, liquid, and solid is reflected by the equilibrium contact angle. The Young [15] equation determines the contact angle for the perfectly flat surface:

$$\gamma_{SG} - \gamma_{SL} - \gamma_{LG} \cos \theta_C = 0. \hspace{1cm} (5)$$

The wettability is significantly affected by the surface roughness and contamination, whereby the Young equation is not globally valid for these surfaces. However, for more or less rugged surface, if it is wetted homogenously, the contact angle can be determined by the Wenzel [16] correlation.

$$\cos(\theta_m) = r \cos(\theta_C), \hspace{1cm} (6)$$

where $r$ is the roughness ratio and is determined as the actual surface area divided by the projected area.

In order to study this phenomenon, the contact angle was measured by the Stalder [17] technique. Figure 3 shows a drop of pure water contact angle with the coated and uncoated surfaces. As shown in this figure, increasing the EPD time or increasing the thickness of the coating, the contact angle was decreased; in other words, the wettability of surfaces increased with SiO$_2$ nanocoating.

3.2. Effect of Sintering. Four plates coated with EPD were sintered in 200, 500, 700, and 800 degrees of centigrade for 5 hours. Sintering, the process of compacting and forming the...
particles by heat or pressure without melting them, is one of the methods of increasing the strength and durability of the nanoparticle coating. As shown in Figure 4, sintering has a significant effect on the contact angle, according to Forrest et al. [18]; the surface is converted to a super-hydrophilic surface by sintering. Table 1 shows the contact angle of various coated plates.

3.3. Coating Thickness. To understand the deposition process, the kinetic deposition of EPD can be used. Initially, when the deposition resistance is not increased, the deposition can be correlated by the linear Hamaker [19] law. Hamaker explained his experimental measurements with linear EPD kinetics and showed that the mass of the deposition linearly depends on the suspension properties and can be written as

$$z = \frac{C \mu Et}{\rho}.$$  \hspace{1cm} (7)

To determine the thickness of SiO\textsubscript{2} nanocoating, the cross section SEM images were captured. In a microscale view shown in Figure 5, the coating is homogenous and smooth. Figure 6 shows the thickness of the deposition with the two percentage of the nanofluid. As shown in this figure, the coating thickness is strongly affected by the particle percentage in the nanofluid. Figure 7 shows the SEM images of electrophoretic deposition at various deposition times, and it can be noted that the mass of deposition increases with time. A comparison between the thicknesses of silica nanocoating versus time, as made in this study, and Charlot et al. [20] experiment trend line are shown in Figure 8(a). In Figure 8(b), the thickness versus concentration is compared with the Hamaker model, equation (7).

In Figure 8(a), the coating experimental results of this study are compared with the experimental data of Charlot et al. [20] at 0.5% concentration. The figure shows that results compare well. However, Figure 8(b) shows that increasing the nanofluid concentration, the Hamaker model predicts the measured coating thickness. This finding indicates that the control of the coating thickness is easier at lower concentrations, and therefore, the EPD and subsequent experimental tests were conducted at lower concentrations.

3.4. Pool Boiling Results. As shown in Table 1, the contact angle of the coated plates was decreased; therefore, the surfaces become hydrophilic and super-hydrophilic. The boiling curve obtained for the bare and SiO\textsubscript{2} coated plates is shown in Figure 9. The significant increment at the superheat temperature of the wall for the hydrophilic plates is conspicuous. The effect of sintering on the boiling of the coated plates is demonstrated in Figure 10. As shown in the figure, the superheat temperature under the nucleate boiling
The regime is decreased at low sintering temperatures (200 and 500°C) and is significantly increased at high sintering temperatures (700 and 800°C). In the course of this study, the effects of the modified surfaces created by the EPD technique on the nucleate boiling heat transfer and the CHF have been investigated.

3.4.1. Heat Transfer Coefficient. As depicted in Figures 9 and 10, the boiling curves of the coated plates tend to shift to higher wall superheat than the clean surface in the nucleate boiling regime. This implies that decreasing the surface contact angle or increasing wettability obtained by coating and then sintering (Table 1), as Wang and Dhir [21] predicted, the nucleate boiling heat transfer occurs at higher surface temperatures, whereas the boiling heat transfer coefficient is decreased.

Figure 11 implies that the nucleate heat transfer coefficient is decreased with various coated surfaces which have lower contact angle than bare surface. The heat transfer coefficient versus heat flux for bare and sintered surfaces (first 40 min EPD coating and then sintered at 200, 500, 700, and 800°C) is shown in Figure 12. However, it can be seen in this figure that for two sintered plates (200 and 500°C), the heat transfer is enhanced compared to coated and unsintered surfaces. Particle adhesion to each other and to heating surface at low sintering temperatures increases without affecting or decreasing the surface porosity. This phenomenon increases the heat transfer coefficient compared to the unsintered surfaces. However, increasing the sintering temperature, the particles surface melting rate increases and the porous layer and the capillary suction both decrease. On the other hand, the reduction of the nucleation site density due to the reduction of the porosity and capillary suctions lead to the reduction of the heat transfer coefficient.

To explain these processes, the widely used Mikic and Rohsenow [22] nucleate boiling heat transfer correlation can be used:

\[
q = \frac{1}{2} \left( \pi k_i \rho c_p \right)^{1/2} f^{1/2} D_p^2 n_a \Delta T_s, \tag{8}
\]

\[
\Delta T_s = \frac{q}{1/2 \left( \pi k_i \rho c_p \right)^{1/2} f^{1/2} D_p^2 n_a} \tag{9}
\]

Table 1: Contact angle and thickness of various coated surfaces.

| Case                       | Contact angle (°C) |
|----------------------------|--------------------|
| Bare plate                 | 88.4               |
| 20 minute with 0.5% nanofluid | 77.4               |
| 60 minute with 0.5% nanofluid | 45.5               |
| Sintered at 200°C          | 37.4               |
| Sintered at 500°C          | 24.6               |
| Sintered at 700°C          | 12.9               |
| Sintered at 800°C          | 5.4                |

Figure 3: Static contact angles of pure water on stainless steel plates for (a) clean surface and (b) 20 minutes and (c) 60 minutes EPD coating.

Figure 4: Static contact angles of pure water on stainless steel plates for (a) 200°C, (b) 500°C, (c) 700°C, and (d) 800°C.
As shown in equation (9), the surface superheat temperature or the heat transfer coefficient depends on the nucleation site density, bubble departure diameter, and bubble departure frequency. Phan et al. [23] observed that the nucleation site density remains almost constant with changing of the contact angle at averaged time period. The departure bubble diameter widely predicted by the Fritz [24] correlation is
Equation (10) relates only the fluid properties and the gravity effects and shows the maximum bubble volume as a function of capillary length and contact angle. The bubble departure is determined by surface tension and buoyancy forces. Hypothetically, the departure occurs when buoyancy overcomes the surface tension, inertia, and drag forces that prevent the detachment of the bubbles.

Phan et al. [23] obtained a correlation that incorporates the Fritz model assumptions with the influence of the energy factor as the contribution of the wetting effects:

\[ d_{bd} = 0.0208\theta_1 \frac{\sigma}{g\Delta\rho} \]  \hspace{1cm} (10)

\[ f(\theta) = \frac{2 + 3\cos\theta - \cos^3\theta}{4} \]  \hspace{1cm} (11)

\[ D_d = 0.626977 f(\theta) \left( \frac{\sigma}{g(\rho_l - \rho_v)} \right)^{1/2} \]  \hspace{1cm} (12)

where \( f(\theta) \) is a volume ratio between a truncated sphere and a complete sphere with equal diameter. Equation (12) shows that the bubble diameter increases with increasing surface wettability. The required energy to activate nucleate sites as well as the growth time \( (t_g) \) and the waiting time \( (t_w) \) is increased for the bigger bubbles. The bubble departure frequency can be written by

Figure 8: Thickness silica nanocoating versus (a) deposition time (dash line is the trend line in Charlot et al. experiment [20]) and (b) concentration (dash line is the Hamaker model [19]).

Figure 9: Pool boiling curves for bare and coated plates.

Figure 10: Pool boiling curves for bare and sintered plates.
In other words, at higher surface wettability, the bubble departure frequency decreases, but the bubble diameter increases. Figure 13 shows the bubble departure frequency and bubble diameter versus the contact angle from Phan et al. [23, 25].

As shown in Figure 13, the bubble departure frequency, between 70 and 90 degrees and lower than 10 degrees, has the highest slip reduction, but the bubble diameter is reduced between 50 and 90, and for lower than 50, it becomes almost a constant value. The wall superheat increment at higher wettability can also be explained using equation (9).

\[ f = \frac{1}{t_g + t_w} \]  

\[ q_n^{CHF} = K\rho_f h_f g \left[ \frac{g\sigma (\rho_f - \rho_g)}{\rho_g^2} \right]^{1/4} \]  

where K depends on the system geometry.

According to the hydrodynamic theory, the CHF is a Taylor-driven process and the vapor jets have the highest Helmholtz instability velocity that arises at the nodes of the Taylor waves. The hydrodynamic theory only depends on the fluids properties and does not consider the surface parameter such as wettability and surface modifications. Figure 9 illustrates that the CHF increases by the nanoparticle deposition on the heated surface; it also shows the limitation of the Zuber model.

### 3.4.2. CHF

Formanyyears, the CHF for departure from the nucleate boiling of pool boiling based on the hydrodynamic instability theory. For many years, the CHF determines the liquid microlayer with the bubble departure frequency affect the outcome of the heat transfer coefficient for hydrophilic and super-hydrophilic surfaces.

\[ f = \frac{1}{t_g + t_w} \]  

\[ q_n^{CHF} = K\rho_f h_f g \left[ \frac{g\sigma (\rho_f - \rho_g)}{\rho_g^2} \right]^{1/4} \]  

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According to the hydrodynamic theory, the CHF is a Taylor-driven process and the vapor jets have the highest Helmholtz instability velocity that arises at the nodes of the Taylor waves. The hydrodynamic theory only depends on the fluids properties and does not consider the surface parameter such as wettability and surface modifications. Figure 9 illustrates that the CHF increases by the nanoparticle deposition on the heated surface; it also shows the limitation of the Zuber model.

In this study, for better understanding, the heat transfer coefficient versus contact angle for various heat fluxes is shown in Figure 14. It can be deduced that the effect of bubble departure frequency on the heat transfer coefficient at the contact angle range of 90°–70° and lower than 30° is dominant, as reported by Phan et al. [23, 25].

For hydrophilic surfaces, the effect of liquid microlayer underneath the bubble is an important parameter that the bigger area of liquid microlayer and lower thickness of microlayer improve the heat transfer coefficient. Therefore, the competition between the area and the thickness of the liquid microlayer with the bubble departure frequency affect the outcome of the heat transfer coefficient for hydrophilic and super-hydrophilic surfaces.
The hot/dry spot theory proposed by Theofanous and Dinh [27] is one of the theories that investigate the surface parameters. According to this theory, CHF develops when rewetting is unable to occur at a growing hot spot. The CHF proposed by this theory is

\[
q_{\text{CHF}}'' = \frac{g \rho_f (\rho_f - \rho_g) \left( \frac{\rho_f^2}{\rho_g^2} \right)^{1/4}}{\kappa - \frac{1}{2}},
\]

where \( \kappa \) is a constant which depends on the surface parameter and is explicitly obtained by Jun et al. [5].

\[
\kappa = \left( 1 - \frac{\sin \theta}{2} \right) \left( \frac{\pi}{2} - \theta \right)^{-(1/2)}.
\]

This correlation is only valid for hydrophilic surfaces.

Another correlation that depends on the surface parameter is the Kandlikar [28] model which can be used for all contact angles.

\[
q_{\text{CHF}}'' = \frac{1}{2 \pi} \frac{(1 + \cos \beta)}{16} \times \left( \frac{2 \pi}{4} (1 + \cos \beta) \cos \phi \right)^{1/2}
\]

\[
\cdot \left[ \frac{\sigma g (\rho_f - \rho_g)}{\rho_f} \right]^{1/4},
\]

where \( \beta \) is the receding contact angle.

Liao et al. [29] represented the Zuber CHF correlation with the effect of orientation and contact angle:

\[
q_{\text{CHF}}'' = \frac{55 - \theta}{100} (0.56) \rho_f h_f \left[ \frac{\sigma g (\rho_f - \rho_g)}{\rho_f^2 \rho_g^2} \right]^{1/4}.
\]

A comparison between the CHFs obtained in this study, the three CHF correlations (15), (17), and (18), and the experimental data [30–32] as a function of the contact angle is shown in Figure 15.

The Zuber model predicts the constant CHF value, but the Kandlikar and the Liao CHF correlations are not constant, and the CHF decreases with the increase in contact angle. These reductions are linear in the Liao correlation and are nonlinear in the Kandlikar correlation. The Kandlikar correlation overpredicts the CHF experimental data. However, the CHF results of this study are closer to those of Hsu and Chen [32].

4. Conclusion

In this investigation, the EPD on thin plates was used to study the critical heat flux behavior in a pool boiling experiment. Because of the negatively charged particles, the particles deposited on the anodic electrode. The main findings of the present investigation can be summarized as follows:

(i) The EPD technique, due to the setup simplicity, shape independency, uniformity of particle coating, and controllability of thickness, can be an appropriate option for modification of the surface and coating on the nuclear fuel cladding.

(ii) The heat transfer coefficient of the nanocoated surfaces decreased from 27 to 11 kW/m²K at 298 kW/m² wall heat flux by nearly 100%.

(iii) Increasing the deposition time and increasing the coating thickness (up to 90 nm for 0.5% nanofluid) increased the CHF.

(iv) The results obtained in this study show that the Kandlikar correlation overpredicts (on average about 30%) the CHF for nanocoated surfaces.

(v) The CHF of the nanocoated surfaces increased from 630 to 1000 kW/m² by 58% and increased from 1000 to 1664 kW/m² by 60% for sintered surfaces.
Nomenclature

C: Concentration (g/cm³)
d: Diameter
E: Electric field (V/cm)
f: Frequency
g: Gravity acceleration (m/s²)
h: Enthalpy (J/kg)
N: Nucleation site density
I: Ampere (A)
q': Heat flux (W/m²)
R: Resistance (Ω)
T: Temperature (°C)
t: Time (s)
V: Voltage (volt)
Z: Thickness (cm).

Greek Letters

ρ: Density (g/cm³)
μ: Electrophoretic mobility (cm²/V·s)
γ: Surface energy
θ: Contact angle (radian)
σ: Surface tension.

Subscripts

B: Bubble
CHF: Critical heat flux
eff: Effective
W: Wall
Sat: Saturate
S: Solid
L: Liquid
G: Gas or vapor.

Data Availability

The data used to support the findings of this study are available from the corresponding author upon request.

Additional Points

Highlights. (i) An anodic electrophoretic deposition (EPD) was used as nanoparticle coating. (ii) The hydrophilic and super-hydrophilic surfaces were established by SiO₂ coating. (iii) Higher strength and CHF were obtained by sintering of the coated surfaces. (iv) The BHT and CHF behaviors of hydrophilic and super-hydrophilic surfaces have been investigated. (v) The EPD technique could be an appropriate option for coating of nuclear fuel claddings.

Conflicts of Interest

The authors declare that they have no conflicts of interest.

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