The transiting system HD 15337: a pair of nearly equal-mass sub-Neptunes on opposite sides of the radius gap

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We report the discovery of two small planets transiting the star HD 15337 (TOI-402, TIC 120896927), a bright ($V = 9$) K1 V dwarf observed by the Transiting Exoplanet Survey Satellite (TESS) in Sectors 3 and 4. We combine the TESS photometry with archival HARPS spectra to confirm the planetary nature of the transit signals and derive the masses of the two transiting planets. With an orbital period of 4.8 days, a mass of $7.63 \pm 0.94 M_\oplus$ and a radius of $1.585 \pm 0.056 R_\oplus$, HD 15337 b joins the growing group of short-period planets known to have a rocky terrestrial composition. HD 15337 c has an orbital period of 17.2 days, a mass of $7.37^{+1.63}_{-1.61} M_\oplus$ and a radius of $2.309^{+0.110}_{-0.103} R_\oplus$, suggesting that the planet might be surrounded by a thick atmospheric envelope. The two planets have virtually the same masses and lie on opposite sides of the radius gap, and are thus an excellent testbed for planet formation and evolution theories. Assuming that HD 15337 c hosts a hydrogen-dominated envelope, we employed a recently developed planet atmospheric evolution algorithm in a Bayesian framework to estimate the history of the high-energy emission of the host star concluding that at 150 Myr it ranged between 1.5 and 93 times that of the current Sun.

Keywords: Planetary systems – Planets and satellites: individual: HD 15337 b – Stars: fundamental parameters – Stars: individual: HD 15337 – Techniques: photometric – Techniques: radial velocities

1. INTRODUCTION

Successfully launched in April 2018, NASA’s Transiting Exoplanet Survey Satellite (TESS) is providing a significant step in understanding the diversity of exoplanets. TESS is performing an all-sky photometric search for planets transiting bright stars ($6 < V < 11$), so that detailed characterizations of the planets and their atmospheres can be performed (Ricker et al. 2015). The survey is broken up into 26 sectors – each sector being observed for ~28 days and consisting of 4 cameras with a combined field-of-view of $24^\circ \times 96^\circ$. Candidate alerts and full-frame images are released every month. As of March 2019, TESS has already announced the discovery of about a dozen transiting planets (see, e.g., Esposito et al. 2018; Gandolfi et al. 2018; Huang et al. 2018; Jones et al. 2018; Nielsen et al. 2018; Quinn et al. 2019; Trifonov et al. 2019).

TESS has already led to the detection of promising systems in terms of planet atmospheric characterization, such as π Men, which is a bright ($V = 5.65$) star hosting a transiting super-Earth with a bulk density consistent with either a primary, hydrogen-dominated atmosphere, or a secondary, probably CO$_2$/H$_2$O-dominated atmosphere (Gandolfi et al. 2018; Huang et al. 2018). The discovery of such systems is central for performing multi-wavelength transmission spectroscopy observations to identify the nature of the atmosphere, thus constraining planetary atmospheric evolution models.

TESS enables also the discovery of multi-planet systems for which both mass and radius can be precisely measured. Since such planets orbit the same star, differences in average density and atmospheric structure among planets belonging to the same system can be ascribed mainly to differences in planetary mass and orbital separation (see, e.g., Guenther et al. 2017; Prieto-Arranz et al. 2018). This greatly simplifies modeling of their past evolution history, thus constraining how these planets formed (Alibert et al. 2005; Alibert & Benz 2017). In this respect, even more significant are multi-planet systems in which two or more planets have similar masses, as differences in radii would most likely be due to the different orbital separations.

In this paper we report the discovery of two small planets transiting the bright ($V = 9$) star HD 15337 (Table 1), a K1 dwarf observed by TESS in Sectors 3 and 4. We combined the TESS photometry with archival HARPS radial velocities (RVs) to confirm the planetary nature of the transit signals and derive the masses of the two planets. The paper is organized as follows. In Sect. 2, we present the TESS photometry and the detection of the transit signals. In Sect. 3, we present the archival HARPS spectra. The properties of the host star are reported in Sect. 4. We present the frequency analysis of the HARPS RVs in Sect. 5 and the data modeling in Sect. 6. Results, discussions, and summary are given in Sect. 7.

2. TESS PHOTOMETRY

HD 15337 (TIC 120896927) was observed by TESS Camera #2 in Sectors 3 and 4 (CCDs #3 and #4, respectively) from 20 September 2018 to 15 November 2018, and will not be observed further during the nominal two-year TESS mission. Photometry was interrupted when the satellite was re-pointed for data downlink, from BJD$_{TDB} = 2458395.4$ to BJD$_{TDB} = 2458396.6$ in Sector 3, and from BJD$_{TDB} = 2458423.5$
Table 1. Main identifiers, coordinates, proper motion, parallax, and optical and infrared magnitudes of HD 15337.

| Parameter | Value       | Source          |
|-----------|-------------|-----------------|
| **Main identifiers**                                      |
| HD        | 15337       |                  |
| HIP       | 11433       | Hipparcos       |
| TIC       | 120896927   | TIC              |
| TOI       | 402         | TESS Alerts     |
| Gaia DR2  | 5068777809824976256 | Gaia DR2     |
| **Equatorial coordinates**                               |
| RA (J2000.0) | 02° 27′ 28.3781″ | Gaia DR2   |
| DEC (J2000.0) | −27° 38′ 06.7417″ | Gaia DR2   |
| **Proper motion and parallax**                           |
| μα, cos δ (mas yr⁻¹) | −73.590 ± 0.057 | Gaia DR2   |
| μδ (mas yr⁻¹)    | −211.614 ± 0.082 | Gaia DR2   |
| Parallax (mas)   | 17.01 ± 0.81   | Gaia DR2   |
| **Magnitudes**                                           |
| B_T         | 10.170 ± 0.027 | Tycho-2
| V_T         | 9.184 ± 0.018  | Tycho-2
| G           | 8.8560 ± 0.0002 | Gaia DR2
| G_HIP       | 9.3194 ± 0.0011 | Gaia DR2
| G_RP        | 8.2708 ± 0.0016 | Gaia DR2
| J            | 7.553 ± 0.019  | 2MASS
| H            | 7.215 ± 0.034  | 2MASS
| K_s         | 7.044 ± 0.018  | 2MASS
| W1(3.35 μm)  | 6.918 ± 0.054  | ALLWISE
| W2(4.6 μm)   | 7.048 ± 0.020  | ALLWISE
| W3(11.6 μm)  | 7.015 ± 0.017  | ALLWISE
| W4(22.1 μm)  | 6.916 ± 0.072  | ALLWISE

**NOTE**—(a) TESS Input Catalog (Stassun et al. 2018b); (b) Gaia Data Release 2 (Gaia Collaboration et al. 2018); (c) Tycho-2 Catalog (Hog et al. 2000); (d) Two-micron All Sky Survey (Cutri et al. 2003); (e) Wide-field Infrared Survey Explorer catalog (Cutri & et al. 2013).

Figure 1. 5′ × 5′ 103aE emulsion image taken in 1955 from the POSS1-E survey, with the SPOC photometric aperture overplotted in blue (TESS pixel size is 21″), and the positions of Gaia DR2 sources (J2015.5) within 2″ of HD 15337 indicated by circles. HD 15337 is in red, nearby sources contributing more than 1% of their flux to the aperture are in orange (see Sect. 2), and other sources are in green. Due to the proper motion of HD 15337, there is a ∼14″ offset between its Gaia position and its position in the image.

There is an additional data gap in Sector 4 from BJD$_{TDB}$ = 2458424.6 to BJD$_{TDB}$ = 2458421.2, which was caused by an interruption in communications between the instrument and spacecraft.

TESS objects of interest (TOIs) are announced publicly via the MIT TESS data alerts web portal. TOIs 402.01 and 402.02 were announced on 16 January 2019 and 31 January 2019, respectively, in association with the HD 15337 photometry. The TESS pixel data and light curves produced by the Science Processing Operations Center (SPOC; Jenkins et al. 2016) at NASA Ames Research Center were subsequently made publicly available via the Mikulski Archive for Space Telescopes (MAST).2 We iteratively searched the SPOC light curves for transit signals using the Box-least-squares algorithm (BLS; Kovács et al. 2002), after fitting and removing stellar variability using a cubic spline with knots every 1.0 days. We recovered two signals corresponding to the TOIs, but no other significant signals were detected. We also tried removing variability using the wavelet-based filter routines VARLET and PHALET, but it did not change the BLS results; we are thus confident that the two signals are robustly detected and are not the result of data artefacts resulting from the choice of variability model or residual instrumental systematic signals. The phase-folded transit signals are shown in Fig. 6.

The SPOC light curves are produced using an optimized aperture, which is shown in Fig. 1. We produced light curves from the TESS pixel data using a series of apertures (Gandolfi et al. 2018; Esposito et al. 2018),

1 https://tess.mit.edu/alerts.

2 https://mast.stsci.edu.
and found that larger apertures than the SPOC aperture shown in Fig. 1 minimized the 6.5 hour combined differential photometric precision (CDPP) noise metric (Christiansen et al. 2012). However, the transit signals recovered from these light curves were slightly less significant, which we attribute to the improvement in light curve quality afforded by the Presearch Data Conditioning (PDC; Smith et al. 2012; Stumpe et al. 2012) pipeline used by the SPOC, which corrects for common-mode systematic noise; for this reason, we opted to analyze the SPOC light curves for the remainder of the analysis in this paper.

To investigate the possibility of diluting flux from stars other than HD 15337, we visually inspected archival imaging and compared Gaia DR2 (Gaia Collaboration 2018) source positions with the SPOC photometric aperture. We used the coordinates of HD 15337 from the TESS Input Catalog3 (TIC; Stassun et al. 2018b) to retrieve Gaia DR2 sources using a search radius of 3′. In archival imaging taken in 1955 from the POSSI survey4, HD 15337 is offset from its current position by ~14′ due to proper motion, but this is not sufficient to completely rule out chance alignment with a background source; however, such an alignment with a bright source is highly unlikely. Assuming the TESS point spread function (PSF) can be approximated by a 2D Gaussian profile with a FWHM of ~25′′, we found that 98.5% of the flux from HD 15337 is within the SPOC aperture. Approximating the TESS bandpass with the Gaia G_RP bandpass, the transit signals from HD 15337 should be diluted by less than 0.01%; HD 15337 is the only star contributing flux to the aperture that is bright enough to be the source of the transit signals. Two other Gaia DR2 sources (5068777809825770112 and 506877745400963584) are also within the SPOC aperture, but they are too faint to yield significant dilution (G_RP ≈ 19 mag). Fig. 1 shows the archival image, along with Gaia DR2 source positions and the SPOC photometric aperture.

3. HARPS SPECTROSCOPIC OBSERVATIONS

HD 15337 was observed between December 2003 and September 2017 UT with the High Accuracy Radial velocity Planet Searcher (HARPS) spectrograph (R ≈ 115,000, Mayor et al. 2003) mounted at the ESO 3.6 m telescope, as part of the observing programs 072.C-0488, 183.C-0972, 192.C-0852, 196.C-1006, and 198.C-0836. We retrieved the publicly available reduced spectra from the ESO archive, along with the cross-correlation function (CCF) and its bisector, computed from the dedicated HARPS pipeline using a K5 numerical mask (Baranne et al. 1996). On June 2015, the HARPS fiber bundle was upgraded and a new set of octagonal fibers, with improved mode-scrambling capabilities, were installed (Lo Curto et al. 2015). To account for the RV offset caused by the instrument refurbishment, we treated the HARPS RVs taken before/after June 2015 as two different data sets. Tables 3 and 4 list the HARPS RVs taken with the old and new fiber bundle, along with the RV uncertainties, the full width at half maximum (FWHM) and bisector span (BIS) of the CCF, the exposure times, and the signal-to-noise ratio (S/N) per pixel at 5500 Å. Time stamps are given in barycentric Julian Date in the barycentric dynamical time (BJD_TDB). We rejected two data points – marked with asterisks in Tables 3 and 4 – because of poor S/N ratio (BJD_TDB = 2455246.519846) or systematics (BJD_TDB = 2457641.794439).

4. STELLAR FUNDAMENTAL PARAMETERS

4.1. Spectroscopic parameters

We co-added the HARPS spectra obtained with the old and new fiber bundle separately to get two combined spectra with S/N per pixel at 5500 Å of 590 (old fiber) and 490 (new fiber). We derived the spectroscopic parameters of HD 15337 from the co-added HARPS spectra using Spectroscopy Made Easy (SME), a spectral analysis tool that calculates synthetic spectra and fits them to high-resolution observed spectra using a χ² minimizing procedure. The analysis was performed with the non-LTE SME version 5.2.2, along with ATLAS 12 one-dimensional model atmospheres (Kurucz 2013).

We estimated a micro-turbulent velocity of $v_{\text{mic}} = 0.80 \pm 0.10$ km s$^{-1}$ from the empirical calibration equations for Sun-like stars from Bruntt et al. (2010). The effective temperature $T_{\text{eff}}$ was measured fitting the wings of the Hα and Hβ lines, as well as the NaI doublet at 5890 and 5896 Å (Fuhrmann et al. 1993; Axer et al. 1994; Fuhrmann et al. 1994, 1997b,a). The surface gravity log $g_*$ was determined from the wings of the Ca i λ 6102, λ 6122, λ 6162Å triplet, and the Ca i λ 6439 Å line, as well as from the Mg i λ 5167, λ 5173, λ 5184 Å triplet. We measured the iron abundance [Fe/H], the macroturbulent velocity $v_{\text{mac}}$, and the projected rotational velocity $v \sin i_*$ by simultaneously fitting the unblended iron lines in the spectral region 5880–6600 Å.

Our analyses applied to the two stacked HARPS spectra provide consistent results, well within the uncertainties. The final adopted values are listed in Table 2. We derived an effective temperature of $T_{\text{eff}} =$
5125±50 K, surface gravity log $g_\star = 4.40±0.10$ (cgs), and an iron abundance relative to solar of [Fe/H] = 0.15 ± 0.08 dex. We also measured a calcium abundance of [Ca/H] = 0.08 ± 0.04 dex and a sodium abundance of [Na/H] = 0.27 ± 0.054 dex. We found a macro-turbulent velocity of $v_{\text{mac}} = 3.0 ± 1.0$ km/s in agreement with the value predicted from the empirical equations of Doyle et al. (2014). The projected rotational velocity was found to be $v \sin i_\star = 1.0 ± 1.0$ km s$^{-1}$.

### 4.2. Stellar mass, radius, age and interstellar extinction

We performed an analysis of the broadband spectral energy distribution (SED) together with the Gaia Data Release 2 (DR2; Gaia Collaboration et al. 2018) parallax in order to determine an empirical measurement of the stellar radius, following the procedures described in Stassun & Torres (2016), Stassun et al. (2017), and Stassun et al. (2018a). We retrieved the $B_\lambda$ and $V_\lambda$ magnitudes from the Tycho-2, the Strömgren $ubvy$ magnitudes from APASS, the $BVgriz$ magnitudes from APASS, the $JHK_s$ magnitudes from 2MASS (Cutri et al. 2003), the W1–W4 magnitudes from ALLWISE (Cutri et al. 2013), and the $G$ magnitude from Gaia DR2 (Gaia Collaboration et al. 2018). Together, the available photometry spans the full stellar SED over the wavelength range 0.35–22 µm (Fig. 2). In addition, we pulled the NUV flux from GALEX in order to assess the level of chromospheric activity, if any.

We performed a fit using Kurucz stellar atmosphere models, with the fitted parameters being the effective temperature $T_{\text{eff}}$ and iron abundance [Fe/H], as well as the interstellar extinction $A_V$, which we restricted to the maximum line-of-sight value from the dust maps of Schlegel et al. (1998). The broadband SED is largely insensitive to the surface gravity (log $g_\star$), thus we simply adopted the value from the initial spectroscopic analysis presented in the previous subsection. The resulting fit is excellent (Figure 2) with a reduced $\chi^2$ of 2.3 (excluding the GALEX NUV flux, which is consistent with a modest level of chromospheric activity). The best fit parameters effective temperature and iron content are $T_{\text{eff}} = 5130 ± 50$ K and [Fe/H] = $0.1^{+0.2}_{-0.1}$ dex, respectively, in excellent agreement with the spectroscopic values (Sect. 4.1). We found that the reddening of HD 15337

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**Figure 2.** Spectral energy distribution (SED) of HD 15337. Red symbols represent the observed photometric measurements, where the horizontal bars represent the effective width of the passband. Blue symbols are the model fluxes from the best-fit Kurucz atmosphere model (black).

**Figure 3.** Hertzsprung-Russell Diagram for HD 15337 based on the observed effective temperature and bolometric luminosity, the latter computed directly from $F_{bol}$ and the Gaia parallax-based distance. Each panel compares the observed properties of the star to evolutionary tracks from the Yonsei-Yale models (Yi et al. 2001; Spada et al. 2013) for different permitted combinations of stellar mass and metallicity. Blue points with labels represent the model ages in Gyr. The central panel represents that case most compatible with all of the available data, including the stellar age of $\approx 5.1$ Gyr as determined from the observed chromospheric activity and stellar rotation period (see the text).
is consistent with zero ($A_v = 0.02 \pm 0.02$ mag), as expected given the relatively short distance to the star (≈45 pc). Integrating the unreddened model SED gives the bolometric flux at Earth of $F_{\text{bol}} = 7.29 \pm 0.08 \times 10^{-9}$ erg s cm$^{-2}$. Taking the $F_{\text{bol}}$ and $T_{\text{eff}}$ together with the Gaia DR2 parallax, adjusted by $+0.08$ mas to account for the systematic offset reported by Stassun & Torres (2018), gives the stellar radius as $R_\star = 0.856 \pm 0.017 R_\odot$. Finally, estimating the stellar mass from the empirical relations of Torres et al. (2010) and a 6% error from the empirical relation itself gives a stellar mass of $M_\star = 0.91 \pm 0.06 M_\odot$.

We can refine the stellar mass estimate by taking advantage of the observed chromospheric activity, which can constrain the age of the star via empirical relations. For example, taking the chromospheric activity indicator, log $R'_{\text{HK}} = -4.916 \pm 0.038$ from Gomes da Silva et al. (2014) and applying the empirical relations of Mamajek & Hillenbrand (2008), gives a predicted age of 5.1 ± 0.8 Gyr. As shown in Fig. 3, according to the Yonsei-Yale stellar evolutionary models (Yi et al. 2001; Spada et al. 2013), this age is most compatible with a stellar mass of $M_\star = 0.90 \pm 0.03 M_\odot$ and [Fe/H] = 0.25, which with the empirically determined stellar radius implies a stellar log $g_\star = 4.52 \pm 0.02$ (cgs), in good agreement with the spectroscopic value of log $g_\star = 4.40 \pm 0.10$ (cgs).

Other combinations of stellar mass and metallicity are compatible with the observed effective temperature and radius (Fig. 3), however they require ages that are incompatible with that predicted by the chromospheric $R'_{\text{HK}}$ emission. Finally, we can further corroborate the activity-based age estimate by also using empirical relations to predict the stellar rotation period from the activity. For example, the empirical relation between $R'_{\text{HK}}$ and rotation period from Mamajek & Hillenbrand (2008) predicts a rotation period for this star of ≈42 days, which is compatible with the observed rotation period derived from the HARPS RVs and activity indicators ($P_{\text{rot}} = 36.5$ days; see the following section).

5. FREQUENCY ANALYSIS OF THE HARPS MEASUREMENTS

We performed a frequency analysis of the HARPS time-series to search for the Doppler reflex motion induced by the two transiting planets discovered by TESS. We accounted for the RV offset between the two different set-ups of the instrument (old and new fiber bundle) using the value of 19.6 m s$^{-1}$ derived from the joint analysis presented in Sect. 6, which is in good agreement with the expected offset for a slowly rotating K1 V star, such as HD 15337 (Lo Curto et al. 2015).

The offset-corrected HARPS RVs are displayed in Fig. 4 (upper panel), along with the FWHM and BIS time-series (middle and lower panel, respectively). The generalized Lomb-Scargle (GLS) periodogram (Zechmeister & Kürster 2009) of the combined RV data shows significant power at frequencies lower than the inverse of the temporal baseline of the HARPS observations. A similar trend is observed in the FWHM obtained with the old fiber bundle (middle panel, blue circles), suggesting that the RV trend might be due to long-term stellar variability (e.g., magnetic cycles).

Alternatively, the RV trend might be induced by a long period orbiting companion, while the long-term variation of the FWHM might be ascribable to the steady instrument de-focusing observed between 2004 and 2015 (Lo Curto et al. 2015).

The upper panel of Fig. 5 shows the GLS periodogram of the combined HARPS RVs, following the subtraction of the best-fitting quadratic trend (cfr. Fig. 4). The peaks with the highest power are found at the orbital frequencies of the two transiting planets ($f_c = 0.058$ c/d and $f_b = 0.210$ c/d), with false-alarm probabilities$^6$ (FAPs) of ≈1% and RV semi-amplitude of about 2.0-2.5 m s$^{-1}$. The periodogram of the RV residuals after subtracting the signal of the outer planet (Fig. 5, sec-

$^5$ We note that the FWHM and BIS offsets between the two instrument set-ups are unknown.

$^6$ The false-alarm probability was derived using the bootstrap method described in (Kuerster et al. 1997).
Figure 5. Generalized Lomb-Scargle periodograms of 1) the combined HARPS RV measurements, following the subtraction of the quadratic trend (first panel); 2) the RV residuals after subtracting the signal of HD 15337 c (second panel); 3) the RV residuals after subtracting the signal of HD 15337 b and c (third panel); the FWHM of the cross-correlation function (fourth panel); the bisector span (BIS) of the cross-correlation function (fifth panel). The dashed horizontal lines mark the false-alarm probability at 0.1, 1 and 5 %. The frequencies of the two transiting planets, as well as of the signal at 36.5 days are marked with vertical arrows.

The transiting multi-planet system HD 15337

On the first panel, shows a significant peak (FAP < 0.1 %) at the frequency of the inner planet. The two peaks have no counterparts in the periodograms of the activity indicators (FWHM and BIS; Fig. 5, fourth and fifth panel), suggesting that the signals are induced by two orbiting planets with periods of 4.8 and 17.2 days. Finally, the GLS periodogram of the RV residuals after subtracting the quadratic trend and the signals of the two planets (Fig. 5, third panel) displays a peak with a FAP < 0.1 % at ~36.5 days, which is also significantly detected in the periodogram of the FWHM (fourth panel). We interpreted the 36.5-day signal as the rotation period of the star. It agrees with the value expected from the $R'_{HK}$.

We combined the activity indicators from the two HARPS fibers by subtracting the best fitting second order polynomials shown in Fig. 4.
activity indicator (Sect. 4.2) and from combining the stellar projected rotational velocity and stellar radius.

6. JOINT ANALYSIS

We performed a joint analysis of the TESS light curve (Sect. 2) and RV measurements (Sect. 3) using the software suite pyaneti, which allows for parameter estimation from posterior distributions calculated using Markov chain Monte Carlo methods.

We extracted ~8 hours of TESS data points centered around each of the transits observed by TESS during Sector 3 and 4. We de-trended the photometric segments with the code exotrending using a second-order polynomial fitted to the out-of-transit data. As described in Sect. 3, we rejected 2 HARPS RVs and used the remaining 85 Doppler measurements, while accounting for an RV offset between the two different HARPS setups.

The RV model includes a linear and a quadratic term, to account for the long-term variation described in Sect. 5, as well as two Keplerians, to account for the Doppler reflex motion induced by HD 15337 b and HD 15337 c. The RV stellar signal at the star’s rotation period was modeled as an additional coherent sine-like curve whose period was constrained with a uniform prior centered at $P_{\text{rot}} = 36.5$ days and having a width of 0.2 days, as derived from the FWHM of the peak detected in the periodogram of the HARPS FWHMs. For the phase and amplitude of the activity signal we adopted uniform priors. While this simple model might not fully reproduce the periodic and quasi-periodic variations induced by evolving active regions carried around by stellar rotation, it has proven to be effective in accounting for the stellar signal of active and moderately active stars (e.g., Pepe et al. 2013; Gandolfi et al. 2017; Barragán et al. 2018). Any variation not properly modeled by the coherent sine-curve, and/or any instrumental noises not included in the nominal RV uncertainties, were accounted for by fitting two RV jitter terms for the two HARPS setups.

We modeled the TESS transit light curves using the limb-darkened quadratic model of Mandel & Agol (2002). For the limb darkening coefficients, we set Gaussian priors using the values derived by Claret (2017) for the TESS pass-band. We imposed conservative error bars of 0.1 on both the linear and the quadratic limb-darkening term. For the eccentricity and argument of periastron we adopted the parametrization proposed by Anderson et al. (2011). A preliminary analysis showed that the transit light curve poorly constrains the scaled semi-major axis $(a/R_*)$. We therefore set a Gaussian prior on $a/R_*$ using Kepler’s third law, the orbital period, and the derived stellar mass and radius (Sect. 4.2).

We imposed uniform priors for the remaining fitted parameters. Details of the fitted parameters and prior ranges are given in Table 2. We used 500 independent Markov chains initialized randomly inside the prior ranges. Once all chains converged, we used the last 5000 iterations and saved the chain states every ten iterations. This approach generates a posterior distribution of 250,000 points for each fitted parameter. Table 2 lists the inferred planetary parameters. They are defined as the median and 68% region of the credible interval of the posterior distributions for each fitted parameter. The transit and RV curves are shown in Fig. 6.

We also experimented with Gaussian Processes (GPs) to model the correlated RV noise associated with stellar activity. Gaussian Processes model stochastic processes with a parametric description of the covariance matrix. GP regression has proven to be successful in modeling the effect of stellar activity for several other exoplanetary systems (see, e.g., Haywood et al. 2014; Grunblatt et al. 2015; López-Morales et al. 2016; Barragán et al. 2018). To this aim, we modified the code pyaneti in order to include a GP algorithm coupled to the MCMC method. We implemented the GP approach proposed by Rajpaul et al. (2015). Briefly, this framework assumes that the star-induced RV variations and activity indicators can be modeled by the same underlying GP and its derivative. This allows the GP to disentangle the RV activity component from the planetary signals.

We assumed that the stellar activity can be modeled by the quasi-periodic kernel described by Rajpaul et al. (2015). We modeled together the HARPS RV, BIS, and FWHM time-series and we treated RV and BIS as being described by the GP and its first derivative, while for FWHM we assumed that it is only described by the GP. The fitted hyper-parameters are then $V_c$, $V_r$, $C_c$, $C_r$, $L_c$, as defined by Rajpaul et al. (2015), to account for the GP amplitudes of the RV, BIS, and FWHM signals, the period of the activity signal $P_{\text{GP}}$, the inverse of the harmonic complexity $\lambda_p$, and the long term evolution timescale $\lambda_c$. We coupled this GP approach with the joint modeling described in the previous section (omitting the extra coherent signal).

As for the planetary signals, we imposed the same priors described in the previous subsection. For the hyper-parameters, we used uniform priors, except for $P_{\text{GP}}$, for which we imposed a Gaussian prior with mean 36.5 d and standard deviation of 0.2 d. We used 250 chains to explore the parameters space. We created the posterior distributions with 500 iterations of converged chains, this generated a posterior distribution with 250,000 points for each parameters.
For planets b and c we derived an RV semi-amplitude of $2.71^{+0.54}_{-0.51} \text{ m s}^{-1}$ and $2.06^{+0.64}_{-0.58} \text{ m s}^{-1}$, respectively, which are in very good agreement with the values reported in Table 2. The other planetary and orbital parameters are also consistent with the values presented in Table 2. For the GP hyper-parameters, we found $V_c = 0.55 \pm 0.23 \text{ m s}^{-1}$, $V_r = 7.0^{+2.7}_{-1.9} \text{ m s}^{-1}$, $B_c = 9.4^{+3.0}_{-2.3} \text{ m s}^{-1}$, $B_r = 6.4^{+20}_{-29} \text{ m s}^{-1}$, $L_c = 5.4 \pm 2.2 \text{ m s}^{-1}$, $P_{GP} = 36.5 \pm 0.2 \text{ d}$, $\lambda_c = 4217^{+624}_{-685} \text{ d}$, and $\lambda_p = 1086^{+501}_{-394} \text{ d}$. The relatively large values of the scale parameters in the GP, i.e. $\lambda_c$ and $\lambda_p$, indicate that the stellar activity behaves like a sinusoidal signal (with slight corrections).

7. DISCUSSION AND CONCLUSIONS

The inner-most transiting planet HD 15337 b ($P_{\text{orb},b} = 4.8 \text{ days}$) has a mass of $M_b = 7.63 \pm 0.94 \text{ M}_\oplus$ and a radius of $R_b = 1.585 \pm 0.056 \text{ R}_\oplus$, yielding a mean density of $\rho_b = 10.5^{+1.8}_{-1.6} \text{ g cm}^{-3}$. Figure 7 displays the position...
of HD 15337 b on the mass-radius diagram compared to
the sub-sample of small transiting planets \( R \leq 4 R_\oplus \)
whose masses and radii have been derived with a pre-
cision better than 25\%. Theoretical models from Zeng
et al. (2016) are displayed with different lines and colors. The
solid red and blue circles mark the position of HD 15337 b
and HD 15337 c, respectively.

For HD 15337, we obtained a mass of \( M_\star = 7.37^{+1.63}_{-1.61} M_\odot \)
and a radius of \( R_\star = 2.309^{+0.110}_{-0.103} R_\odot \), yielding a mean
density of \( \rho_\star = 3.3^{+0.9}_{-0.8} \) g cm\(^{-3}\). Therefore, HD 15337 b
and c have virtually the same mass, but the radius of
HD 15337 c is \(~1.5\) larger than the radius of HD 15337 b.
The lower bulk density of HD 15337 c suggests that the
planet is likely composed by a rocky core surrounded
either by a considerable amount of water, or by a light,
hydrogen-dominated envelope. In the first case, the
amount of water and high planetary equilibrium tem-
perature would imply the presence of a steam atmo-
sphere, which would be strongly hydrogen dominated
in its upper part as a consequence of water dissociation
and of the low mass of hydrogen. It is therefore plausible
to assume that HD 15337 c hosts a hydrogen-dominated
atmosphere, at least in its upper part.

As in other systems hosting two close-in sub-Neptune-
mass planets (e.g., HD 3167 Gandolfi et al. 2017), the
radii of HD 15337 b and c lie on opposite sides of the ra-
dius gap (Fulton et al. 2017; Van Eylen et al. 2018),
with the closer-in planet having a higher bulk den-
sity, similarly to what observed in other close-in sys-
tems with measured planetary masses (e.g., K2-109,
HD 3167, GJ 9827 Guenther et al. 2017; Gandolfi et al.
2017; Prieto-Arranz et al. 2018). This gap is believed
to be caused by atmospheric escape (Owen & Wu 2017;
Jin & Mordasini 2018), which is stronger for closer-in
planets. Within this context, HD 15337 b has probably
lost its primary, hydrogen-dominated atmosphere, and
now hosts a secondary atmosphere possibly resulting
from out-gassing of a solidifying magma ocean, while
HD 15337 c is likely to still partly retain the primordial
hydrogen-dominated envelope. This is consistent with
Van Eylen et al. (2018), who measured the location and
gle of the radius gap as a function of orbital period
and matched it to models suggesting a homogeneous ter-
restrial core composition.

The present-day properties of the atmospheres of
HD 15337 b and c are therefore intimately related to the
amount of high-energy (X-ray and EUV; <91.2 nm) stel-
lar radiation received since the dispersal of the proto-
planetary nebula, and thus also to the stellar rotation
history. The evolution of the stellar rotation rate does
not follow a unique path because stars of the same mass
and metallicity can have significantly different rotation
rates up to about 1 Gyr (e.g., Mamajek & Hillenbrand
2008; Johnstone et al. 2015; Tu et al. 2015). For older
stars, it is therefore not possible to infer their past high-
energy emission from their measured stellar properties.
Starting from the assumption that HD 15337 c hosts a
hydrogen-dominated atmosphere, we derive the history
of the stellar rotation and high-energy emission by mod-
eling the atmospheric evolution of HD 15337 c. To this end, we employ the atmospheric evolution algorithm described by Kubyshkina et al. (2018) and further developed by Kubyshkina et al. (2019, submitted), which is based on a Bayesian approach fitting the currently observed planetary radius, combining the planetary evolution model with the open-source MCMC algorithm of Cubillos et al. (2017). The planetary atmospheric evolution model, system parameters (i.e., planetary mass, planetary radius, orbital separation, current stellar rotation period, stellar age, stellar mass) are then used to compute the posterior distribution for the stellar rotation rate at any given age via MCMC. We assumed Gaussian priors determined by the measured system parameters and their uncertainties.

Figure 8 shows the obtained posterior distribution for the rotation period HD 15337 at an age of 150 Myr in comparison with the distribution derived from measurements of open cluster stars of the same age (Johnstone et al. 2015). Our results indicate that HD 15337, when it was young, was likely to be a moderate rotator, with a high-energy emission at 150 Myr ranging between 1.5 and 93 times the current solar emission, further excluding that the star was a very fast/slow rotator.

The position of HD 15337 c in the mass-radius diagram (Fig. 7) indicates that the planet may be hosting a massive hydrogen-dominated envelope or a smaller secondary atmosphere. As primary atmospheres are easily subject to escape, knowing the current composition of the envelope of HD 15337 c would provide a strong constrain on atmospheric evolution models. In this respect, this planet is similar to π Men c (Gandolfi et al. 2018; Huang et al. 2018); furthermore, as for π Men, the close distance to the system and brightness of the host star would enable high-quality transmission spectroscopy observations spanning from far-ultraviolet to infrared wavelengths. Of particular interest would be far-ultraviolet observations covering H1, C1i, and O1 resonance lines, that could be obtained with the Hubble Space Telescope and would ideally probe the nature of the atmosphere, thus constraining atmospheric escape and evolution models.

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**Facilities:** TESS, HARPS

**Software:** VARLET (Grziwa & Pätzold 2016), PHALET (Grziwa & Pätzold 2016), SME (Valenti & Piskunov 1990; Valenti & Fischer 2005; Piskunov & Valenti 2017), exotrending (Barragán & Gandolfi 2017), pyaneti (Barragán et al. 2019)

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Table 2. HD 15337 system parameters.

| Parameter                                      | Prior(s) | Derived value |
|------------------------------------------------|----------|---------------|
| **Stellar parameters**                        |          |               |
| Star mass $M_*$ ($M_\odot$)                   |          | 0.90 ± 0.03   |
| Star radius $R_*$ ($R_\odot$)                 |          | 0.856 ± 0.017 |
| Effective Temperature $T_{\text{eff}}$ (K)     |          | 5125 ± 50     |
| Surface gravity (b) log $g_*$ (cgs)            |          | 4.52 ± 0.02   |
| Surface gravity (c) log $g_*$ (cgs)            |          | 4.40 ± 0.10   |
| Iron abundance [Fe/H] (dex)                    |          | 0.15 ± 0.10   |
| Projected rotational velocity $v \sin i_*$ (km s$^{-1}$) |          | 1.0 ± 1.0     |
| Age (Gyr)                                      |          | 5.1 ± 0.8     |
| Interstellar extinction $A_*$                  |          | 0.02 ± 0.02   |
| **Model parameters of HD 15337 b**             |          |               |
| Orbital period $P_{\text{orb}, b}$ (days)      | $U[4.73, 4.78]$ | 4.756216 ± 0.000081 |
| Transit epoch $T_{0, b}$ (BJD$_{\text{DB}}$−2 450000) | $U[8411.35, 8411.55]$ | 8411.46165 ± 0.00057 |
| Scaled semi-major axis $a_b/R_*$               | $N[13.11, 0.17]$ | 13.44 ± 0.30   |
| Planet-to-star radius ratio $R_b/R_*$          | $U[0, 0.1]$ | 0.01698 ± 0.00050 |
| Impact parameter $b_0$                         | $U[0, 1]$ | 0.30 ± 0.15   |
| $\sqrt{\cos i_b} \sin \omega_{e, b}$         | $U[-1, 1]$ | 0.13 ± 0.12   |
| $\sqrt{\cos i_b} \cos \omega_{e, b}$         | $U[-1, 1]$ | 0.20 ± 0.21   |
| Radial velocity semi-amplitude variation $K_b$ (m s$^{-1}$) | $U[0, 10]$ | 3.13 ± 0.39   |
| **Model parameters of HD 15337 c**             |          |               |
| Orbital period $P_{\text{orb}, c}$ (days)      | $U[17.08, 17.28]$ | 17.17753 ± 0.00077 |
| Transit epoch $T_{0, c}$ (BJD$_{\text{DB}}$−2 450000) | $U[8414.4, 8414.7]$ | 8414.55162 ± 5.38150 |
| Scaled semi-major axis $a_c/R_*$               | $N[13.11, 0.17]$ | 31.68 ± 0.73   |
| Planet-to-star radius ratio $R_c/R_*$          | $U[0, 0.1]$ | 0.0247 ± 0.0010 |
| Impact parameter $b_c$                         | $U[0, 1]$ | 0.88 ± 0.03   |
| $\sqrt{\cos i_c} \sin \omega_{e, c}$         | $U[-1, 1]$ | −0.07 ± 0.29  |
| $\sqrt{\cos i_c} \cos \omega_{e, c}$         | $U[-1, 1]$ | −0.02 ± 0.17  |
| Radial velocity semi-amplitude variation $K_c$ (m s$^{-1}$) | $U[0, 10]$ | 1.98 ± 0.43   |
| **Additional model parameters**                |          |               |
| Parameteralized limb-darkening coefficient $q_1$ | $N[0.43, 0.1]$ | 0.37 ± 0.10   |
| Parameteralized limb-darkening coefficient $q_2$ | $N[0.19, 0.1]$ | 0.18 ± 0.10   |
| Systemic velocity $\gamma_{\text{HS1}}$ (km s$^{-1}$) | $U[-4.0, -3.6]$ | −3.8172 ± 0.0024 |
| Systemic velocity $\gamma_{\text{HS2}}$ (km s$^{-1}$) | $U[-4.0, -3.6]$ | −3.7976 ± 0.0011 |
| RV jitter term $\sigma_{\text{HS1}}$ (m s$^{-1}$) | $U[0, 100]$ | 1.99 ± 0.24   |
| RV jitter term $\sigma_{\text{HS2}}$ (m s$^{-1}$) | $U[0, 100]$ | 2.65 ± 0.43   |
| Stellar rotation period ($P_{\text{rot}}$) days | $U[36.4, 36.6]$ | 36.539 ± 0.024 |
| Linear RV term $m$ s$^{-1}$ d$^{-1}$           | $U[-100, 100]$ | −0.0056 ± 0.0015 |
| Quadratic RV term $m$ s$^{-1}$ d$^{-2}$        | $U[-100, 100]$ | −0.000013 ± 0.000002 |
| **Derived parameters of HD 15337 b**           |          |               |
| Planet mass $M_b$ ($M_\oplus$)                |          | 7.63 ± 0.94   |
| Planet radius $R_b$ ($R_\oplus$)              |          | 1.585 ± 0.056 |
| Planet mean density $\rho_b$ (g cm$^{-3}$)     |          | 10.5$^{+1.8}_{-1.6}$ |
| Semi-major axis of the planetary orbit $a_c$ (AU) |          | 0.0535 ± 0.0016 |
| Orbit eccentricity $e_c$                       |          | 0.09 ± 0.06   |
| Orbit inclination $i_c$ (deg)                  |          | 88.7$^{+0.8}_{-0.6}$ |
| Equilibrium temperature (c) $T_{\text{eq}, c}$ (K) |          | 989 ± 15     |
| Transit duration $\tau_{14, c}$ (hours)        |          | 2.518 ± 0.036 |
| **Derived parameters of HD 15337 c**           |          |               |
| Planet mass $M_c$ ($M_\oplus$)                |          | 7.37$^{+1.63}_{-1.61}$ |
| Planet radius $R_c$ ($R_\oplus$)              |          | 2.309$^{+0.110}_{-0.103}$ |
| Planet mean density $\rho_c$ (g cm$^{-3}$)     |          | 3.3$^{+0.9}_{-0.8}$ |
| Semi-major axis of the planetary orbit $a_c$ (AU) |          | 0.1261 ± 0.0038 |
| Orbit eccentricity $e_c$                       |          | 0.08$^{+0.10}_{-0.06}$ |
| Orbit inclination $i_c$ (deg)                  |          | 88.4 ± 0.1    |
| Equilibrium temperature (c) $T_{\text{eq}, c}$ (K) |          | 644 ± 10     |
| Transit duration $\tau_{14, c}$ (hours)        |          | 2.229$^{+0.070}_{-0.060}$ |

**Note:**  (a) $U[a, b]$ refers to uniform priors between $a$ and $b$, and $F[a]$ to a fixed $a$ value; (b) from spectroscopy and isochrones; (c) from spectroscopy; (d) assuming zero albedo and uniform redistribution of heat.
Table 3. HARPS RV measurements of HD 15337 acquired with the old fiber bundle.

| BJD$^a_{DB}$ | RV  | $\pm \sigma$ | BIS  | FWHM  | $T_{exp}$ | S/N $^b$ |
|--------------|-----|--------------|------|--------|-----------|---------|
| -2450000     | -3.8208 | 0.0010 | 0.0013 | 6.1330 | 900 | 70.8 |
| -3.8260      | 0.0008 | 0.0015 | 6.1353 | 900 | 93.7 |
| -3.8234      | 0.0007 | 0.0010 | 6.1471 | 900 | 101.7 |
| -3.8155      | 0.0006 | 0.0059 | 6.1571 | 900 | 108.1 |
| -3.8111      | 0.0007 | 0.0037 | 6.1531 | 900 | 97.7 |
| -3.8185      | 0.0006 | 0.0049 | 6.1547 | 900 | 101.4 |
| -3.8170      | 0.0005 | 0.0034 | 6.1481 | 900 | 125.8 |
| -3.8203      | 0.0006 | 0.0063 | 6.1631 | 900 | 102.6 |
| -3.8128      | 0.0006 | 0.0100 | 6.1779 | 900 | 108.1 |
| -3.8127      | 0.0007 | 0.0133 | 6.1767 | 900 | 98.8 |
| -3.8132      | 0.0010 | 0.0150 | 6.1661 | 900 | 68.4 |
| -3.8115      | 0.0011 | 0.0053 | 6.1684 | 900 | 62.2 |
| -3.8140      | 0.0008 | 0.0015 | 6.1633 | 1200 | 79.4 |
| -3.8101      | 0.0006 | 0.0061 | 6.1753 | 900 | 112.5 |
| -3.8117      | 0.0007 | 0.0065 | 6.1745 | 900 | 100.7 |
| -3.8107      | 0.0006 | 0.0078 | 6.1681 | 900 | 117.0 |
| -3.8088      | 0.0008 | 0.0102 | 6.1717 | 900 | 85.6 |
| -3.8139      | 0.0006 | 0.0078 | 6.1734 | 900 | 103.1 |
| -3.8116      | 0.0006 | 0.0062 | 6.1635 | 900 | 113.6 |
| -3.8169      | 0.0006 | 0.0088 | 6.1583 | 900 | 113.4 |
| -3.8165      | 0.0009 | 0.0028 | 6.1577 | 900 | 78.1 |
| -3.8191      | 0.0016 | -0.0011 | 6.1644 | 900 | 47.0 |
| -3.8105      | 0.0005 | 0.0036 | 6.1476 | 900 | 120.8 |
| -3.8164      | 0.0007 | 0.0068 | 6.1558 | 900 | 93.6 |
| -3.8185      | 0.0008 | 0.0056 | 6.1655 | 800 | 85.2 |
| -3.8159      | 0.0008 | 0.0071 | 6.1630 | 900 | 84.5 |
| -3.8123      | 0.0008 | 0.0012 | 6.1615 | 900 | 79.2 |
| -3.8063      | 0.0007 | 0.0049 | 6.1668 | 900 | 91.2 |
| -3.8165      | 0.0009 | 0.0080 | 6.1621 | 900 | 76.6 |
| -3.8147      | 0.0009 | 0.0105 | 6.1507 | 900 | 73.4 |
| -3.8138      | 0.0007 | 0.0023 | 6.1464 | 900 | 89.9 |
| -3.8110      | 0.0006 | 0.0041 | 6.1508 | 900 | 116.6 |
| -3.8188      | 0.0008 | 0.0028 | 6.1553 | 900 | 85.7 |
| -3.8125      | 0.0005 | 0.0052 | 6.1507 | 900 | 126.2 |
| -3.8104      | 0.0007 | 0.0066 | 6.1632 | 900 | 94.7 |
| -3.8095      | 0.0006 | 0.0115 | 6.1656 | 900 | 102.7 |
| -3.8120      | 0.0006 | 0.0095 | 6.1629 | 900 | 111.4 |
| -3.8122      | 0.0006 | 0.0038 | 6.1587 | 900 | 104.4 |
| -3.8099      | 0.0006 | 0.0008 | 6.1533 | 900 | 104.2 |
| -3.8161      | 0.0005 | 0.0008 | 6.1593 | 900 | 120.8 |
| -3.8150      | 0.0007 | 0.0091 | 6.1510 | 900 | 98.7 |
| -3.8161      | 0.0007 | -0.0025 | 6.1456 | 900 | 94.4 |
| -3.8120      | 0.0007 | 0.0087 | 6.1559 | 900 | 97.8 |
| -3.8169      | 0.0360 | 0.1136 | 6.3367 | 5 | 3.8 |
| -3.8094      | 0.0007 | 0.0047 | 6.1576 | 900 | 94.9 |
| -3.8098      | 0.0007 | -0.0005 | 6.1539 | 1200 | 92.8 |
| -3.8089      | 0.0010 | 0.0033 | 6.1612 | 900 | 69.8 |
| -3.8073      | 0.0008 | 0.0024 | 6.1609 | 900 | 92.2 |
| -3.8119      | 0.0013 | 0.0010 | 6.1674 | 900 | 58.4 |
| -3.8060      | 0.0008 | 0.0040 | 6.1663 | 900 | 90.9 |
| -3.8091      | 0.0009 | 0.0001 | 6.1772 | 900 | 83.9 |
| -3.8147      | 0.0008 | 0.0041 | 6.1679 | 900 | 88.9 |

Note—(a) Barycentric Julian dates are given in barycentric dynamical time; (b) S/N per pixel at 550 nm; (*) outlier not included in the analysis.
Table 4. HARPS RV measurements of HD 15337 acquired with the new fiber bundle.

| BJD\(_{TDB}\) | RV \(\pm\sigma\) | BIS | FWHM | \(T_{exp}\) | S/N \(^b\) |
|---------------|-----------------|-----|------|---------|-------|
| 7291.826359   | -3.7848 0.0008 | 0.0273 | 6.2102 | 900 | 88.1   |
| 7292.799439   | -3.7889 0.0007 | 0.0246 | 6.2141 | 900 | 95.6   |
| 7299.843584   | -3.7988 0.0007 | 0.0229 | 6.1968 | 900 | 106.2  |
| 7303.879389   | -3.7973 0.0013 | 0.0184 | 6.1985 | 900 | 60.3   |
| 7353.698805   | -3.7959 0.0008 | 0.0209 | 6.2081 | 900 | 97.6   |
| 7357.682828   | -3.7898 0.0006 | 0.0239 | 6.2073 | 900 | 127.1  |
| 7373.685080   | -3.7889 0.0007 | 0.0207 | 6.2064 | 900 | 99.5   |
| 7395.644348   | -3.7931 0.0011 | 0.0253 | 6.2093 | 900 | 68.5   |
| 7399.617257   | -3.7938 0.0008 | 0.0213 | 6.1981 | 900 | 90.1   |
| 7418.584180   | -3.7951 0.0007 | 0.0267 | 6.1927 | 900 | 108.5  |
| 7422.589672   | -3.7969 0.0008 | 0.0211 | 6.1964 | 900 | 103.2  |
| 7427.538475   | -3.7932 0.0008 | 0.0197 | 6.1990 | 900 | 94.5   |
| 7429.539077   | -3.7883 0.0006 | 0.0222 | 6.2054 | 900 | 129.3  |
| 7584.927622   | -3.7909 0.0007 | 0.0218 | 6.2010 | 900 | 97.2   |
| 7613.935104   | -3.8009 0.0007 | 0.0214 | 6.1960 | 900 | 104.8  |
| *7641.794439* | -1.6179 0.0009 | 0.0220 | 6.1919 | 900 | 81.4   |
| 7642.837840   | -3.8015 0.0009 | 0.0196 | 6.1934 | 900 | 80.2   |
| 7643.808690   | -3.8000 0.0008 | 0.0189 | 6.1928 | 900 | 89.1   |
| 7644.862657   | -3.7982 0.0009 | 0.0202 | 6.1907 | 900 | 79.1   |
| 7647.923051   | -3.7972 0.0007 | 0.0181 | 6.1888 | 900 | 110.7  |
| 7649.726575   | -3.7895 0.0010 | 0.0205 | 6.1981 | 900 | 72.2   |
| 7650.752860   | -3.7912 0.0006 | 0.0207 | 6.2017 | 900 | 115.2  |
| 7652.744706   | -3.7889 0.0008 | 0.0222 | 6.1982 | 900 | 91.5   |
| 7656.751475   | -3.7956 0.0008 | 0.0254 | 6.2022 | 900 | 87.1   |
| 7658.854305   | -3.7910 0.0005 | 0.0196 | 6.1970 | 900 | 142.1  |
| 7660.797461   | -3.7952 0.0008 | 0.0195 | 6.1985 | 900 | 85.9   |
| 7661.831637   | -3.7959 0.0007 | 0.0116 | 6.2990 | 900 | 107.0  |
| 7971.834240   | -3.7981 0.0009 | 0.0219 | 6.1895 | 900 | 79.6   |
| 7993.916286   | -3.8000 0.0009 | 0.0246 | 6.2071 | 900 | 84.7   |
| 7994.887034   | -3.7979 0.0011 | 0.0250 | 6.2032 | 900 | 70.0   |
| 7996.829372   | -3.7884 0.0012 | 0.0185 | 6.2004 | 900 | 64.2   |
| 7996.923718   | -3.7902 0.0009 | 0.0238 | 6.1989 | 1500 | 89.6   |
| 7998.866474   | -3.7986 0.0010 | 0.0216 | 6.2004 | 900 | 72.2   |
| 8001.872927   | -3.7930 0.0008 | 0.0188 | 6.1873 | 900 | 99.5   |
| 8002.895001   | -3.7975 0.0009 | 0.0216 | 6.1880 | 900 | 82.8   |

Note—(a) Barycentric Julian dates are given in barycentric dynamical time; (b) S/N per pixel at 550 nm; (*) outlier not included in the analysis.
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