Spectra–Structure Correlations in Isotopomers of Ethanol (CX₃CX₂OX; X = H, D): Combined Near-Infrared and Anharmonic Computational Study

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Academic Editor: Fabrizio Santoro
Received: 8 May 2019; Accepted: 8 June 2019; Published: 11 June 2019

Abstract: The effect of isotopic substitution on near-infrared (NIR) spectra has not been studied in detail. With an exception of few major bands, it is difficult to follow the spectral changes due to complexity of NIR spectra. Recent progress in anharmonic quantum mechanical calculations allows for accurate reconstruction of NIR spectra. Taking this opportunity, we carried out a systematic study of NIR spectra of six isotopomers of ethanol (CX₃CX₂OX; X = H, D). Besides, we calculated the theoretical spectra of two other isotopomers (CH₃CD₂OD and CD₃CH₂OD) for which the experimental spectra are not available. The anharmonic calculations were based on generalized vibrational second-order perturbation theory (GVPT2) at DFT and MP2 levels with several basis sets. We compared the accuracy and efficiency of various computational methods. It appears that the best results were obtained with B2PLYP-GD3BJ/def2-TZVP//CPCM approach. Our simulations included the first and second overtones, as well as binary and ternary combinations bands. This way, we reliably reproduced even minor bands in the spectra of diluted samples (0.1 M in CCl₄). On this basis, the effect of isotopic substitution on NIR spectra of ethanol was accurately reproduced and comprehensively explained.

Keywords: near-infrared spectroscopy; ethanol; anharmonic quantum mechanical calculations; isotopic substitution; overtones; combinations bands

1. Introduction

Near-infrared (NIR) spectra are appreciably more complex and difficult for interpretation than IR or Raman spectra [1–6]. This results from a large number of strongly overlapping overtones and combination bands, numerous resonances between different modes and anharmonicity of vibrations [1–6]. Interpretation of vibrational bands has been aided by studies of a series of similar compounds (including isotopomers), or by reconstruction of the spectra by using quantum mechanical calculations [5]. The former way may provide highly speculative assignments, while the latter method has limitations that prevent their common use in NIR spectroscopy. From the point of view of applied spectroscopy, there exists an essential difference in the applicability of quantum mechanical calculation of mid-infrared (MIR) [7–9] and NIR spectra. A simplistic and computationally inexpensive harmonic approximation fails to predict the overtones and combination modes [10]. Because of a considerable computational cost of anharmonic calculations, simulations of NIR spectra are rare. Nevertheless, in the literature one can find examples of application of different approaches used for calculation of overtones.
and combination bands, including variational approaches, which are expensive but useful in selected cases [11,12]. However, the studies using a vibrational self-consistent field (VSCF) approach [13,14] and its refined variants—e.g., PT2-VSCF—are more common [15–17]. Recently, a number of theoretical reconstructions of NIR spectra by means of efficient vibrational second-order perturbation (VPT2) method have been reported [18]; e.g., carboxylic acids [19], fatty acids [20–22], aminoacids [23], nucleobases [24], nitriles [25], azines [26], phenols [27,28], and alcohols [29,30]. Considerable efforts have been undertaken in order to develop anharmonic approaches applicable to even larger molecular systems [31–34]. On the other hand, meticulous probing of vibrational potential capable of yielding nearly-exact results is also available [35–38]. Recent advances in this field include the development of multi-dimensional approaches that provide complete information on mode couplings in linear triatomic molecules [39].

The isotopic effect appears to be helpful for the analysis of NIR spectra [10]. By shifting a part of overlapped contributions, one can reduce their complexity and reveal individual bands. Time-resolved NIR spectroscopy of deuterated alcohols has also been successfully used for elucidating the diffusion coefficients [40]. In our previous work, the effect of isotopic substitution on NIR spectra of methanol has been accurately reproduced by anharmonic calculations [41]. In particular, we were able to predict the vibrational contributions from non-uniformly substituted CX₃OX (X = H, D) species, which are not available from the experiment [41]. Further studies on molecules more complex than methanol are still necessary. A reasonable progress in this field is expected by examination of ethanol and all its isotopomers [42,43]. Ethanol has eight major isotopomers resulting from deuteration (CX₃CX₂OX; X = H, D), as compared to four in methanol. Moreover, the internal rotation around C-O(H) bond leads to rotational isomerism (gauche, trans) in molecules of ethanol, which adds additional origin of spectral variability in NIR spectra [42,43].

To enable detailed examination of the impact of various effects on NIR spectra of ethanol isotopomers, at first it is necessary to perform reliable theoretical reproduction of NIR spectra for eight isotopomers of ethanol (CX₃CX₂OX; X = H, D). We are interested in accurate reproduction of subtle effects observed in NIR spectra. Therefore, a combination of several electronic methods underlying VPT2 vibrational analysis will be useful for establishing the best approach capable of reproducing fine features in NIR spectra. The determination of electronic structure underlying the geometry optimization and harmonic analysis will be based on Møller–Plesset second-order perturbation (MP2) and density functional (DFT) theories. The efficiency and accuracy of reproduction of NIR bands by MP2 and DFT with single-hybrid B3LYP and double-hybrid B2PLYP density functionals will be overviewed. MP2 and DFT calculations included basis sets of increasing quality (6-31G(d,p), SNST, def2-TZVP, and aug-cc-pVTZ). Moreover, the impact of solvent cavity model will be evaluated. The anharmonic vibrational analysis will be carried out by means of generalized VPT2 (GVPT2). In our previous studies on methanol, it has been demonstrated that the relative contributions from the second overtones and ternary combinations are different for various isotopomers [41]. This work will provide detailed information on contributions from different vibrational modes and the trends observed with increasing of the alkyl chain length in going from methanol to ethanol. In addition, we will elucidate the accuracy of prediction of the three quanta transitions in NIR spectra.

2. Results and Discussion

2.1. Accuracy of Reproduction of NIR Spectra by Selected Approaches

Anharmonic calculations are significantly more challenging compared to harmonic approximation [1–4]. This holds even for efficient anharmonic approaches based on VPT2 method. At the same time, the higher quanta transitions are more prone to inaccuracies than the fundamental ones [44]. An insufficient accuracy of the ground state geometry and potential energy surface may easily propagate into inaccuracy of prediction in VPT2 calculation step. Thus, the theoretical prediction of NIR
spectra is usually a compromise between the cost and accuracy. Effects like isotopic substitution [41] and conformational isomerism [29,30] may further complicate the vibrational analysis of NIR spectra.

One of the aims of this work was assessment of the efficiency of several combinations of the electronic theory methods and basis sets (Table 1). In addition, we examined the effect of the solvent model on the anharmonic vibrational energies. An efficient single-hybrid B3LYP functional is commonly used tool for spectroscopic studies [10]. Empirical correction for dispersion has been introduced to overcome one of the major weaknesses of DFT method [45]. In some cases, this approach markedly improves the robustness of calculation of primary parameters (i.e., energy). Therefore, recent literature suggests employing empirical correction for dispersion in DFT calculations [46]. Our previous studies have shown that, even for small and isolated molecules, this correction is advantageous in spectra modeling [29]. For molecules in solution an approximation of the solvent cavity often improves the quality of the simulated NIR spectra [28]. However, in the present work this advantage is less important (Table 1). We observed an improvement of RMSE from 45 to 35 cm\(^{-1}\) for \(\text{CH}_2\text{CH}_2\text{OH}\), and from 27 to 24 cm\(^{-1}\) for \(\text{CH}_3\text{CH}_2\text{OD}\). However, considering small additional cost of CPCM (ca. 10% of total CPU time in the case of GVPT2/B3LYP-GD3BJ/6-31G(d,p) calculations), it is advisable to include this correction step in the calculations.

Switching from B3LYP to B2PLYP density functional with a small basis set (6-31G(d,p)) leads to an increase in the RMSE value (Table 1). However, B2PLYP method overestimated the band positions in a systematic manner. In contrast, B3LYP approach provides irregular results. Some of the band positions (\(\delta\text{CH}\)) are blue-shifted, while the others (\(\nu\text{OX}\) and \(\nu\text{CX}; X = \text{H, D}\)) are red-shifted. Thus, more uniform band shift from B2PLYP method (Figures 1 and 2) resulted in better interpretability of NIR spectra as compared with B3LYP results. It has a peculiar effect in NIR spectra of aliphatic alcohols, as it reduces RMSE of NIR band positions, particularly for the \(\nu\text{OX} + \delta\text{CH}\), and \(\nu\text{CX} + \delta\text{CH}\) combination bands. However, this apparent gain does not improve the true interpretability of the spectra. It is likely that simulated NIR spectra of larger molecules may suffer even more due to binary combinations involving the stretching and deformation of the C-H and O-H vibrations. Similar inconsistency of B3LYP as compared to B2PLYP method has been noted before [29,41]. Due to electron correlation being computed effectively at the MP2 level, it is commonly accepted that B2PLYP functional requires larger basis sets. B2PLYP coupled with large def2-TZVP basis set noticeably improves the quality of simulated NIR spectra (Figures 1 and 2). This improvement is particularly evident in the reproduction of minor bands originating from the three quanta transitions (Figures 3 and 4). This effect is nicely illustrated by reduction of RMSE from 72–83 cm\(^{-1}\) for B2PLYP/6-31G(d,p) to 18–19 cm\(^{-1}\) for B2PLYP/def2-TZVP. SNST basis set, less complex than def2-TZVP but still of triple-\(\zeta\) quality, leads to worse results. An exception was observed for the prominent doublet from the \(\nu\text{CH}\) combination bands (at ca. 4400 cm\(^{-1}\)), where B2PLYP/SNST calculations reproduced the peak shapes more resembling the experimental ones. However, the position of this doublet was also overestimated by this method. Hence, B2PLYP/def2-TZVP method appears to be the better tool for reliable reconstruction of NIR spectra. A similar conclusion was obtained for butyl alcohols [30].

In contrast, MP2 method does not appear to be particularly useful for anharmonic calculations of NIR spectra of ethanol isotopomers due to significant redshift of the \(\nu\text{CX}\) frequencies (Figures 1 and 2). It is an interesting observation, as the tendency of MP2 to describe incorrectly repulsive forces is known in the literature [47]. In this work, however, an insufficient basis set (6-31G(d,p)) may strongly deviate the results. Here, MP2 method seems to be more sensitive to the effect of a small basis set than DFT-B2PLYP method. As expected, an application of a larger basis set, e.g., aug-cc-pVTZ, improves the accuracy of calculations. However, these results are not as good as those obtained from B2PLYP/def2-TZVP method. Moreover, this improvement is accompanied by a substantial increase of computing time (by ca. 4 times). Nevertheless, selected spectral ranges (\(\nu\text{OH} + \delta\text{CH} \approx 5050–4800\) cm\(^{-1}\) and \(\nu\text{CH} + \delta\text{CH}\) doublet \(\approx 4400\) cm\(^{-1}\)) were better reproduced by MP2/aug-cc-pVTZ computations. Therefore, MP2 approach may be recommended as a reference method in selected cases. Our results demonstrate that different computational methods achieve different accuracy for particular regions of
NIR spectra, and these regions do not overlap. Thus, comparison of the spectra simulated by different methods (e.g., resulting from VPT2 calculations at DFT or MP2 levels) appears to be the best way for reliable interpretation of the experimental spectra.

Particular attention should be paid to $2\nu_{\text{OH/OD}}$ band, which is the most characteristic peak for alcohols and the other important compounds like, e.g., phenols [27], terpenes [28], and polyphenols [48]. This peak is very sensitive to the chemical environment and inter- and intra-molecular interactions, and is frequently used for studies of the structure and physicochemical properties [49–53]. Hence, its proper theoretical reproduction is of essential importance. As shown in Figures 1–4, and Figures S1–S6 in SM, most of the methods did not reproduce correctly the shape of this band. The experimental band from $2\nu_{\text{OH/OD}}$ vibration reveals a slight asymmetry. This asymmetry is a result of convolution of two components due to trans and gauche conformers. The lower-frequency gauche component has also the lower intensity. Among isotopomers of ethanol, this feature was reproduced correctly only by B2PLYP/def2-TZVP method. MP2/6-31G(d,p) and MP2/aug-cc-pVTZ methods predicted correct shape of the $2\nu_{\text{OH/OD}}$ peak only for CH$_3$CD$_2$OH and CD$_3$CD$_2$OD. As can be seen (Table 1), the peak position was overestimated by all used methods, but B2PLYP calculations give the best agreement.

In comparison with other modes, large amplitude motions (LAMs)—e.g., torsion modes and hindered rotations—are more difficult for accurate description in harmonic approximation and also by anharmonic approaches that probe the potential curve relatively shallow (e.g., VPT2) [54]. We did not find any evidence that these low-frequency modes influence NIR bands directly (i.e., their overtone and combination modes do not appear in NIR region). However, a NIR spectrum provides some insights on LAMs as well. In our case, the shape of the $2\nu_{\text{OH/OD}}$ band is an indirect probe of the accuracy of prediction of the low-frequency modes. The shape of this band results from two components due to gauche and trans rotational conformers. Unreliable theoretical abundances of these forms would result in biased relative intensities of the $2\nu_{\text{OH/OD}}$ components (Table S1). Gibbs free energies may be affected by erroneous LAMs, which would propagate into incorrect relative abundances of gauche and trans conformers. Because it is an isolated band of strong intensity, the simulated $2\nu_{\text{OH/OD}}$ may be used to assess the reliability of prediction of LAMs and the related Gibbs free energies. This kind of error would manifest itself as a distorted shape of simulated $2\nu_{\text{OH/OD}}$ band. Above effect can be seen for some of the methods used in this study, e.g., for B3LYP (B3LYP-GD3BJ/6-31G(d,p); B3LYP-GD3BJ/6-31G(d,p)/CPCM; B3LYP-GD3BJ/ST/ST; and B2PLYP coupled with an insufficient basis set (B2PLYP-GD3BJ/6-31G(d,p)/CPCM.). However, the methods which yielded the most accurate spectra in the other regions (B2PLYP/def2-TZVP; MP2/6-31G(d,p); MP2/aug-cc-pVTZ) also reproduced $2\nu_{\text{OH/OD}}$ peak accurately (Figures 1 and 2). Therefore, we conclude that the LAMs of ethanol and its derivatives were determined adequately by MP2 method. B2PLYP method also provides correct results, but it is more sensitive to the selection of a basis set. On the other hand, B3LYP tends to falsify the Gibbs free energies corrected by anharmonic ZPE. Further studies are needed to determine, whether this effect occurs because of an unreliable description of LAMs. On the other hand, inaccuracy of the $2\nu_{\text{OH/OD}}$ frequencies prediction by VPT2 may also be considered as another contributing factor, as we have evidenced such occurrence in the case of the conformers of cyclohexanol [27]. Note that B3LYP functional coupled with a relatively simple basis set yields reasonable reproduction of NIR spectra and correctly predicts the effects of isotopic substitution at a relatively modest computational expense (Figures 1–4 and Figures S1–S6 in SM). However, a tendency to over- and underestimate the position and intensity of some bands may be unfavorable for the reliable interpretation of theoretical NIR spectra. For exploration of more subtle effects, B2PLYP functional seems to be more suitable. In the present study of isotopic substitution and the other effects (e.g., rotational isomerism) on NIR spectra of ethanol, we used B2PLYP/def2-TZVP method with additions of GD3BJ and CPCM.
Figure 1. NIR spectra of CH$_3$CH$_2$OH calculated with GVPT2 method at different levels of electronic theory; (a) B3LYP-GD3BJ/6-31G(d,p); (b) B3LYP-GD3BJ/6-31G(d,p)//CPCM; (c) B2PLYP-GD3BJ/6-31G(d,p)//CPCM; (d) MP2/6-31G(d,p)//CPCM; (e) B3LYP-GD3BJ/SNST//CPCM; (f) B2PLYP-GD3BJ/def2-TZVP//CPCM; (g) MP2/aug-cc-pVTZ//CPCM; (exp.) Experimental spectrum of CH$_3$CH$_2$OH in CCl$_4$ (0.1 M).
Figure 2. NIR spectra of CH$_3$CH$_2$OD calculated with GVPT2 method at different levels of electronic theory; (a) B3LYP-GD3BJ/6-31G(d,p); (b) B3LYP-GD3BJ/6-31G(d,p)/CPCM; (c) B2PLYP-GD3BJ/6-31G(d,p)/CPCM; (d) MP2/6-31G(d,p)/CPCM; (e) B3LYP-GD3BJ/SNST/CPCM; (f) B2PLYP-GD3BJ/def2-TZVP/CPCM; (g) MP2/aug-cc-pVTZ/CPCM; (exp.) Experimental spectrum of CH$_3$CH$_2$OD in CCl$_4$ (0.1 M).
Figure 3. Contributions from minor bands in NIR spectra of CH$_3$CH$_2$OH calculated with GVPT2 method at different levels of electronic theory; (a) B3LYP-GD3BJ/6-31G(d,p); (b) B3LYP-GD3BJ/6-31G(d,p)//CPCM; (c) B2PLYP-GD3BJ/6-31G(d,p)//CPCM; (d) MP2/6-31G(d,p)//CPCM; (e) B3LYP-GD3BJ/SNST//CPCM; (f) B2PLYP-GD3BJ/def2-TZVP//CPCM; (g) MP2/aug-cc-pVTZ//CPCM; (exp.) Experimental spectrum of CH$_3$CH$_2$OH in CCl$_4$ (0.1 M).
Figure 4. Contributions from minor bands in NIR spectra of CH$_3$CH$_2$OD calculated with GVPT2 method at different levels of electronic theory; (a) B3LYP-GD3BJ/6-31G(d,p); (b) B3LYP-GD3BJ/6-31G(d,p)/CPCM; (c) B2PLYP-GD3BJ/6-31G(d,p)/CPCM; (d) MP2/6-31G(d,p)/CPCM; (e) B3LYP-GD3BJ/SNST/CPCM; (f) B2PLYP-GD3BJ/def2-TZVP/CPCM; (g) MP2/aug-cc-pVTZ/CPCM; (exp.) Experimental spectrum of CH$_3$CH$_2$OD in CCl$_4$ (0.1 M).
Table 1. Positions of selected NIR bands (in cm$^{-1}$) from GVPT2 anharmonic vibrational analysis in CH$_3$CH$_2$OH and CH$_3$CH$_2$OD at different levels of electronic theory and corresponding RMSE values.

| Assignment | Exp. | Calculated | Diff. | RMSE |
|------------|------|------------|-------|-------|
| $\nu_2$OH | 7099 | 7157 58 | 7125 26 | 26 | 7081 –18 | 2743 144 | 7134 35 | 7061 –6 | 2706 –3 |
| $\nu_3$CH$_2$ + $\nu_1$OH | 6520.4 | 6578 57.6 | 6513 –7.4 | 6536 15.6 | 6654 133.6 | 6576 55.6 | 6477 –43.4 | 6472 –48.4 |
| $\nu_2$as'CH$_3$ + $\nu_1$OH | 5886.1 | 5980 93.9 | 5889 2.9 | 5922 35.9 | 6111 224.9 | 5970 83.9 | 5871 –15.1 | 5892 5.9 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 5765.7 | 5897 131.3 | 5790 24.3 | 5864 98.3 | 5957 191.3 | 5898 132.3 | 5797 31.3 | 5812 46.3 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 5665.1 | 5799 93.9 | 5681 15.9 | 5708 42.9 | 5812 146.9 | 5722 56.9 | 5611 –54.1 | 5768 102.9 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 5013.8 | 5026 12.2 | 5029 15.2 | 5012 –1.8 | 5114 100.2 | 5049 35.2 | 5040 26.2 | 5019 5.2 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 4954.2 | 4978 23.8 | 4979 24.8 | 4959 4.8 | 5003 48.8 | 5009 54.8 | 5008 53.8 | 5003 48.8 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 4873 | 4881 8 | 4868 –5 | 4877 4 | 4958 85 | 4926 53 | 4874 1 | 4883 10 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 4394.8 | 4448 53.2 | 4366 –28.8 | 4443 48.2 | 4592 197.2 | 4473 78.2 | 4424 29.2 | 4436 41.2 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 4333.5 | 4395 61.5 | 4331 –2.5 | 4364 30.5 | 4502 168.5 | 4416 82.5 | 4353 19.5 | 4357 23.5 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 4331.8 | 4392 60.2 | 4366 32.2 | 4364 32.2 | 4500 168.2 | 4415 83.2 | 4349 17.2 | 4355 23.2 |
| $\delta$as'CH$_3$ + $\nu_1$OH | 4054.1 | 4057 2.9 | 4037 –17.1 | 4020 –34.1 | 4122 67.9 | 4036 –18.1 | 4068 13.9 | 4058 3.9 |

| RMSE | CH$_3$CH$_2$OH | 70.0 | RMSE | 18.1 | RMSE | 40.8 | RMSE | 153.1 | RMSE | 72.2 | RMSE | 35.0 | RMSE | 44.6 |

| RMSE | CH$_3$CH$_2$OD | 80.6 | RMSE | 18.6 | RMSE | 50.0 | RMSE | 179.4 | RMSE | 83.3 | RMSE | 23.6 | RMSE | 27.3 |
2.2. Origins of NIR Bands of \( \text{CX}_3\text{CX}_2\text{OX} \) \( X = \text{H}, \text{D} \)

The simulations of NIR spectra of ethanol isotopomers in \( \text{CCl}_4 \) solutions by GVP T2 anharmonic method at B2PLYP-GD3BJ/def2-TZVP/CPCM level accurately reproduced most of the experimental bands (Figures 5 and 6). On this basis, we performed detailed and reliable band assignments (Tables 2–9). The consistency of these assignments was positively verified by comparison with the experimental spectra of six isotopomers. High accuracy of simulations allows to analyze the theoretical spectra of \( \text{CH}_3\text{CD}_2\text{OD} \) and \( \text{CD}_3\text{CH}_2\text{OD} \) (Figure 6C,D and Tables 8 and 9) which are not available commercially. All assignments were supported by an analysis of the potential energy distributions (PEDs; Tables S2–S9).

NIR spectra of ethanol isotopomers mainly consist of the combinations of stretching and bending OX and CX (\( X = \text{H}, \text{D} \)) modes (Figures 1–6 and Tables 2–9). The region below \( 5500 \text{ cm}^{-1} \) for \( \text{CH}_3\text{CH}_2\text{OH} \) is almost entirely contributed by the combination bands, while absorption from the overtones dominates above \( 5500 \text{ cm}^{-1} \). NIR spectrum of ethanol may be roughly divided into four regions, but only two of them contain meaningful contributions from overtones. These regions are contributed mainly by vibrations from: (1) \( 2\nu \text{OX} \); (2) \( 2\nu \text{CX} + \nu \text{CX} + \nu \text{CX} \); (3) \( \nu \text{OX} + \delta \text{CX} \); (4) \( \nu \text{CX} + \delta \text{CX} \); \( X = \text{H}, \text{D} \). The other combination bands like \( \nu \text{OX} + \nu \text{CX} \), and \( 2\nu \text{CX} + \nu \text{CX} \) have low intensity. Isotopic substitution introduces significant band shift, strongly affecting the appearance of NIR spectra (Figures 5 and 6 and Figures S1–S6). It should be noted, that the region of \( 5700–5400 \text{ cm}^{-1} \) for ethanol containing \( \text{CH}_3 \) and \( \text{CH}_2 \) groups is strongly affected by the anharmonic effects. This effect is well seen for \( \text{CH}_3\text{CH}_2\text{OH} \) (Figure 1) and \( \text{CH}_3\text{CH}_2\text{OD} \) (Figure 2). The most meaningful contributions in this region originate from \( \nu \text{CH} + \nu \text{CH} \left( \nu_{\text{as}}\text{CH}_2 + \nu_\text{s}\text{CH}_2 \right) \), \( 2\delta \text{CH} + \nu \text{CH} \), and \( 2\nu_{\text{as}}\text{CH}_2 \) vibrations as well.

One can notice the overestimated intensities of the \( 2\nu \text{CH} \) and \( \nu \text{CH} + \nu \text{CH} \) bands appearing in the \( 6000–5500 \text{ cm}^{-1} \) region (Figures 5 and 6 and Figures S1–S6). The magnitude of this effect varies between the different methods; however, it is present in all cases. A similar overestimation we have observed for butyl alcohols [30]. At present, we are unable to explain the reasons for these overestimations. Unexpectedly, B2PLYP functional (regardless of basis set; 6-31G(d,p), SNST, and def2-TZVP yielded similar results) significantly overestimates the frequencies of \( 2\delta \text{CH} + \delta \text{CH} \) transitions, shifting them to the \( 5500 \text{ cm}^{-1} \) region. In contrast, the most of other transitions in NIR region is accurately reproduced by this approach. In the case of the \( 2\delta \text{CH} + \delta \text{CH} \) modes, large positive anharmonic constants appeared in GVP T2 vibrational analysis. Consequently, positions of the corresponding bands were predicted far from a simple combination of the harmonic frequencies. This shift has not been observed for the remaining approaches. Presently, the reason of this behavior is not clear. Because of very low intensity of \( 2\delta \text{CH} + \delta \text{CH} \) bands, these erroneous predictions do not provide meaningful contributions to NIR spectra. However, this occurrence demonstrates the need for using more than one method during examination of the fine spectral effects.

The deuteration of the OH group leads to a noticeable shift of the \( \nu \text{OD} + \delta \text{CH} \) band. In contrast, the other bands do not shift meaningfully, as can be easily seen from comparison of \( \text{CH}_3\text{CH}_2\text{OH} \) and \( \text{CH}_3\text{CH}_2\text{OD} \) spectra (Figure 5A,B). In particular, the absorption from the \( \nu \text{CH} + \delta \text{CH} \) in the \( 4600–4000 \text{ cm}^{-1} \) region remains unaffected. This region can be used to monitor the isotopic substitutions of the \( \text{CH}_3 \) and \( \text{CH}_2 \) groups, as it leads to highly specific spectral changes. Obviously, simultaneous deuteration of both groups implies more significant changes. However, the most interesting effects result from the selective substitution of one of these groups. The presence of the \( \text{CH}_3 \) group gives rise to a prominent doublet near \( 4395 \) and \( 4330 \text{ cm}^{-1} \). This doublet has a complex structure resulting from overlapping of the contributions from the \( \text{CH}_2 \) (Figure S7 in Supplementary Material), leading to a broadening of the high-frequency wing of the doublet. As expected, this contribution is not present in the spectrum of \( \text{CH}_3\text{CD}_2\text{OH} \) (Figure S7B). On the other hand, the isotopic substitution of the \( \text{CH}_3 \) reveals a part of the overlapping contributions, as observed more clearly in the second derivative spectrum of \( \text{CD}_3\text{CH}_2\text{OH} \) (Figure S7C). This effect is well seen in the calculated spectra (Figure S7C).

The higher frequency NIR region \( (>7000 \text{ cm}^{-1}) \) is also very sensitive to the isotopic effect. A weak absorption from the higher order overtones and combination bands creates difficulties in the analysis
of this region. The deuteration of the OH group significantly reduces the number of the bands as a result of red-shift of the combination bands. The simulated spectra confirmed the high isotopic purity of the samples, except of CD$_2$CD$_2$OD which shows the $2\nu$OH peak near 7100 cm$^{-1}$ in the experimental spectrum (Figure 6B). This is in contrast to previously studied methanol, in which various non-uniform substitutions have been identified [41]. Contrary to -CX bonds, the H or D atoms in -OX bonds are labile, therefore, the OD group tends to exchange into the OH even by exposition of the deuterated alcohol to air. Since this band has a high absorptivity, therefore even small impurities due to the OH appear in NIR spectrum as a clear band at 7100 cm$^{-1}$. In contrast, no -CH bands are observed in the spectrum of CD$_3$CD$_2$OD (Figure 6B). NIR spectroscopy is particularly sensitive and selective for the isotopic effect, although the theoretical calculations are necessary for proper spectra interpretation. The spectral manifestations of the OH group in OD derivatives are obscured by ternary combinations from the CH vibrations that appear in the same region. For example, the $\delta_{\text{as}}$CH$_3$ + $\nu_5$CH$_2$ + $\nu_{\text{as}}$CH$_2$ bands in CH$_3$CH$_2$OD (Figure 5B), and $\delta_{\text{scis}}$CH$_2$ + $\nu_5$CH$_2$ + $\nu_{\text{as}}$CH$_2$ bands in CD$_3$CH$_2$OD (Figure 6D) are observed.

One can speculate that the isotopic substitution and conformational isomerism lead to convoluted spectral changes. This phenomenon will be a subject of our next paper (in preparation).

**Figure 5.** Band assignments in NIR spectra of deuterated ethanols based on GVPT2/B2PLYP-GD3BJ/def2-TZVP/PCPM calculations. (A) CH$_3$CH$_2$OH; (B) CH$_3$CH$_2$OD; (C) CH$_3$CD$_2$OH; (D) CD$_3$CH$_2$OH. Band numbering corresponds to that presented in Tables 2–5.
Figure 6. Band assignments in NIR spectra of deuterated ethanols based on GVPT2/B2PLYP-GD3BJ//def2-TZVP//CPCM calculations. (A) CD3CD2OH; (B) CD2CD2OD; (C) CH3CD2OD; (D) CD3CH2OD. Band numbering corresponds to that presented in Tables 6–9.

Table 2. Band assignments in NIR spectra of CH3CH2OH based on GVPT2/B2PLYP-GD3BJ//def2-TZVP//CPCM calculations. Band numbering corresponds to that presented in Figure 5A.

| Peak Number | νExp | νCalc | Assignment (Major Contribution) |
|-------------|------|-------|---------------------------------|
| 1           | 8718.0 | 8739  | 2νasCH2 + νasCH3                 |
| 2           | 8430.0 | 8526  | 3νasCH2                          |
| 3           | 8329.0 | 8416  | 2νs(CH2) + νasCH3 + νas(CH3)     |
| 4           | 8131.0 | 8329  | [δasCH3, δasCH3] + νasCH3 + νas(CH3) |
| 5           | 7400–7300 | 7400–7300 | δs(CH2) + νasCH2 + νas(CH3)      |
| 6           | 7300–7200 | 7300–7200 | δπ(CH3) + νsCH3 + νas(CH3)       |
| 7           | 7099.0 | 7125  | νasCH3 + νOH                     |
| 8           | 6610.0 | 6609  | 2δasCH2 + νOH                    |
| 9           | 6565.0 | 6540  | νCH2 + νOH                        |
| 10          | 6520.4 | 6513  | 2δw(out)CH2 + νOH                |
| 11          | 6331.0 | 6314  | δasCOH + δas(wag)CH2 + νOH       |
| 12          | 6271.8 | 6275  | [νCC, δas(bip)COH] + νasCH3 + νas(CH3) |
Table 3. Band assignments in NIR spectra of CH\(_2\)CH\(_2\)OD based on GVPT2/B2PLYP-GD3BJ/def2-TZVP//CPMC calculations. Band numbering corresponds to that presented in Figure 5B.

| Peak Number | \(\nu_{\text{Exp}}\) | \(\nu_{\text{Calc}}\) | Assignment (Major Contribution) |
|-------------|-------------------|-------------------|-------------------------------|
| 1           | 8428.0            | 8334              | \(2\nu_{\text{as}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_2\) |
| 2           | 7777.5            | 7796              | \(3\nu\text{OD}\)             |
| 3           | 7260.0            | 7227              | \(\delta_{\text{twist}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_3 + \nu_{\text{as}}\text{CH}_3\) |
| 4           | 7099.1            | 7112              | \(\delta_{\text{as}}\text{CH}_3 + \nu_{\text{as}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_2\) |
| 5           | 6133.4            | 6200              | \(\tau\text{CC} + \nu_{\text{as}}\text{CH}_3 + \nu_{\text{as}}\text{CH}_3\) |
| 6           | 5935.0            | 5946              | \(2\nu_{\text{as}}\text{CH}_3\) |
| 7           | 5885.2            | 5895              | \(2\nu_{\text{as}}\text{CH}_3\) |
| 8           | 5850.0            | 5847              | \([\delta_{\text{as}}\text{CH}_3, \delta_{\text{as}}'\text{CH}_3] + \delta_{\text{sciss}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_2\) |
| 9           | 5765.6            | 5788              | \(2\nu_{\text{as}}\text{CH}_2\) |
| 10          | 5665.7            | 5669              | \(2\nu_{\text{as}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_2\) |
| 11          | 5564.3            | 5559              | \(\nu_{\text{OD}} + \nu_{\text{CH}_2}\) |
| 12          | 5494.3            | 5498              | \(\nu_{\text{as}}\text{CH}_2 + \delta_{\text{twist}}\text{CH}_2 + \delta_{\text{wagg}}\text{CH}_2\) |
| 13          | 5277.1            | 5289              | \(2\nu\text{OD}\)             |
| 14          | 4947.0            | 4963              | \([\delta_{\text{rock}}\text{CH}_2, \delta_{\text{rock}}'\text{CH}_2] + \delta_{\text{twist}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_2\) |
| 15          | 4873.0            | 4846              | \([\delta_{\text{rock}}\text{CH}_2, \delta_{\text{rock}}'\text{CH}_2] + [\delta_{\text{rock}}\text{CH}_2, \delta_{\text{rock}}'\text{CH}_2] + \nu_{\text{as}}\text{CH}_2\) |
| 16          | 4720.8            | 4717              | \([\nu_{\text{CO}}, \delta_{\text{rock}}\text{CH}_3, \delta_{\text{ip}}\text{COD}] + [\delta_{\text{rock}}\text{CH}_3, \delta_{\text{ip}}\text{COD}, \delta_{\text{sciss}}\text{CH}_2\text{CO}] + \nu\text{OD}\) |
| 17          | 4393.7            | 4397              | \(\delta_{\text{sciss}}\text{CH}_2 + \nu_{\text{as}}\text{CH}_2\) |
| 18          | 4331.8            | 4364              | \([\delta_{\text{as}}\text{CH}_3, \delta_{\text{wagg}}\text{CH}_2] + \nu_{\text{as}}\text{CH}_2\) |
| 19          | 4253.4            | 4275              | \(\delta_{\text{twist}}\text{CH}_2 + \nu_{\text{as}}'\text{CH}_3\) |
| 20          | 4155.1            | 4127              | \([\delta_{\text{rock}}'\text{CH}_3, \delta_{\text{ip}}\text{COD}, \delta_{\text{sciss}}\text{CH}_2\text{CO}] + \nu_{\text{as}}\text{CH}_3\) |
| 21          | 4105.0            | 4063              | \([\delta_{\text{rock}}'\text{CH}_3, \delta_{\text{ip}}\text{COD}, \delta_{\text{sciss}}\text{CH}_2\text{CO}] + \nu_{\text{as}}\text{CH}_3\) |
| 22          | 4054.1            | 4037              | \([\delta_{\text{rock}}'\text{CH}_2, \delta_{\text{rock}}'\text{CH}_3] + \delta_{\text{CH}}\) |
Table 4. Band assignments in NIR spectra of CH$_3$CD$_2$OH based on GVPT2/B2PLYP-GD3BJ/def2-TZVP//CPM calculations. Band numbering corresponds to that presented in Figure 5C.

| Peak Number | $v_{\text{exp}}$ | $v_{\text{calc}}$ | Assignment (Major Contribution) |
|-------------|------------------|-------------------|---------------------------------|
| 1           | 8717.0           | 8788             | $3v_{as}$CH$_3$                 |
| 2           | 8434.1           | 8728             | $v_{as}$CH$_3 + v_{as}$CH$_3 + v_{as}$'CH$_3$ |
| 3           | 8337.0           | 8665             | $2v_{as}$CH$_3 + v_{as}$CH$_3$ |
| 4           | 7345.0           | 7324             | $[\delta_{\text{rock}} \text{CD}, \delta_{\text{wag}} \text{CD}] + v_{as}$CH$_3 + vOH          |
| 5           | 7251.0           | 7255             | $[\delta_{as} \text{CH}_3, \delta_{as} \text{CH}_3] + v_{as}$CH$_3 + v_{as}$CH$_3$ |
| 6           | 7098.2           | 7126             | 2vOH                           |
| 7           | 6590.4           | 6618             | $v_{as}$'CH$_3 + vOH            |
| 8           | 6205.8           | 6257             | 2$\delta_{ip}$COH + vOH        |
| 9           | 6158.0           | 6198             | 2$\delta_{ip}$COH + vOH        |
| 10          | 5929.9           | 5943             | 2$v_{as}$CH$_3$                 |
| 11          | 5895.4           | 5892             | 2$v_{as}$CH$_3$                 |
| 12          | 5828.0           | 5844             | $v_{as}$CH$_3 + v_{as}$'CH$_3$  |
| 13          | 5763.8           | 5744             | $v_{as}$CD$_2 + vOH            |
| 14          | 5669.1           | 5595             | $v_{as}$CH$_3 + 2v_{as}$CH$_3$  |
| 15          | 5449.0           | 5496             | $\delta_{\text{rock}}$CH$_3 + \delta_{ip}$CH$_3 + v_{as}$'CH$_3$ |
| 16          | 5439.4           | 5449             | $\delta_{\text{sciss}}$CD$_2$CO + 2$\delta_{ip}$CH$_3$ |
| 17          | 5282.0           | 5309             | $v_{as}$CD$_2 + 2[\delta_{as}$CH$_3, \delta_{as}$CH$_3]$|
| 18          | 5070.1           | 5094             | $v_{as}$CD$_2 + v_{as}$CH$_3$   |
| 19          | 5017.0           | 5018             | $[\tau \text{CC}, \delta_{\text{oop}} \text{COH}] + 2[\delta_{as}$CH$_3, \delta_{as}$'CH$_3]$ |
| 20          | 4927.0           | 4954             | $\delta_{ip}$COH + vOH         |
| 21          | 4898.3           | 4930             | $\delta_{ip}$COH + vOH         |
| 22          | 4792.1           | 4806             | $[\nu \text{CO}, \delta_{\text{wag}}$CD$_2] + vOH |
| 23          | 4737.2           | 4744             | $\delta_{\text{sciss}}$CD$_2 + vOH$ |
| 24          | 4650.2           | 4641             | $[\tau \text{CC}, \delta_{\text{rock}} \text{CH}_3] + vOH$ |
| 25          | 4591.7           | 4600             | $[\nu \text{CO}, \delta_{\text{wag}}$CD$_2] + vOH$ |
| 26          | 4513.3           | 4525             | $\delta_{ip}$COH + vOH         |
| 27          | 4404.6           | 4437             | $[\delta_{as}$CH$_3, \delta_{as}$'CH$_3] + v_{as}$CH$_3$ |
| 28          | 4338.0           | 4373             | $\delta_{ip}$CH$_3 + v_{as}$'CH$_3$ |
| 29          | 4275.7           | 4285             | $\delta_{ip}$COH + v_{as}$CH$_3$ |
| 30          | 4238.1           | 4218             | $v_{as}$CD$_2 + v_{as}$CD$_2$   |
| 31          | 4129.6           | 4147             | $[\nu \text{CO}, \delta_{\text{wag}}$CD$_2] + v_{as}$'CH$_3$ |
| 32          | 4079.7           | 4117             | $\delta_{\text{rock}}$CH$_3 + v_{as}$CH$_3$ |
| 33          | 4056.2           | 4052             | $\delta_{\text{rock}}$CH$_3 + v_{as}$CH$_3, \delta_{\text{sciss}}$CD$_2$CO + vOH |

Table 5. Band assignments in NIR spectra of CD$_3$CH$_2$OH based on GVPT2/B2PLYP-GD3BJ/def2-TZVP//CPM calculations. Band numbering corresponds to that presented in Figure 5D.

| Peak Number | $v_{\text{exp}}$ | $v_{\text{calc}}$ | Assignment (Major Contribution) |
|-------------|------------------|-------------------|---------------------------------|
| 1           | 8405.0           | 8511             | $3v_{as}$CH$_2$                 |
| 2           | 8307.6           | 8323             | $2v_{as}$CH$_2 + v_{as}$CH$_2$  |
| 3           | 7098.5           | 7124 (f)         | 2vOH                           |
| 4           | 6841.1           | 6897             | $[\nu \text{CO}, \delta_{as}$'CD$_2] + v_{as}$CH$_2 + v_{as}$CH$_2$ |
| 5           | 6517.6           | 6510             | $v_{as}$CH$_2 + vOH$            |
| 6           | 6324.0           | 6261             | $[\delta_{ip}$COH, $\delta_{\text{twist}}$CH$_2$] + $\delta_{\text{wag}}$CH$_2 + vOH$ |
| 7           | 6268.0           | 6111             | $[\delta_{\text{rock}}$CD$_3$, $\nu$CC] + $v_{as}$CD$_3 + vOH$ |
| 8           | 6070.4           | 6059             | $2[\delta_{ip}$COH, $\delta_{\text{twist}}$H$_2$] + $\nu$OH |
| 9           | 5966.3           | 5966             | $[\delta_{ip}$CD$_3$, $\nu$CC] + $\delta_{\text{wag}}$CH$_2 + vOH$ |
| 10          | 5839.1           | 5825             | 2$v_{as}$CH$_2$                 |
| 11          | 5772.1           | 5793             | 2$v_{as}$CH$_2$                 |
| 12          | 5628.0           | 5686             | 2$v_{as}$CH$_2$                 |
| 13          | 5533.0           | 5641             | 2$v_{as}$CH$_2$                 |
| 14          | 5427.9           | 5419             | $2\delta_{\text{twist}}$CH$_2 + v_{as}$CH$_2$ |
| 15          | 5358.0           | 5367             | $[\nu \text{CO}, \delta_{as}$CD$_2] + v_{as}$CD$_3 + v_{as}$'CD$_3$ |
| 16          | 5286.8           | 5287             | $\delta_{as}$CD$_3 + \delta_{\text{twist}}$CH$_2 + v_{as}$CH$_2$ |
### Table 5. Cont.

| Peak Number | ν<sub>exp</sub> | ν<sub>calc</sub> | Assignment (Major Contribution) |
|-------------|----------------|----------------|---------------------------------|
| 17          | 5190.0         | 5188           | 2[δ<sub>CD3</sub> νCC] + ν<sub>a</sub>CH<sub>2</sub> |
| 18          | 5102.7         | 5084           | δ<sub>oop</sub>COH + 2δwaggCH<sub>2</sub> |
| 19          | 5007.1         | 5028           | δwaggCH<sub>2</sub> + νOH |
| 20          | 4955.8         | 4987           | [δ<sub>twist</sub>CH<sub>2</sub>, δ<sub>p</sub>COH, δwaggCH<sub>2</sub>] + νOH |
| 21          | 4853.8         | 4853           | [δ<sub>p</sub>COH, δ<sub>twist</sub>CH<sub>2</sub>] + νOH |
| 22          | 4764.7         | 4768           | [δ<sub>CD3</sub>, νCC] + νOH |
| 23          | 4676.7         | 4676           | νCO + νOH |
| 24          | 4558.4         | 4511           | δ<sub>a</sub>CD<sub>3</sub> + [δ<sub>twist</sub>CH<sub>2</sub>, δ<sub>p</sub>COH] + ν<sub>a</sub>′′′CD<sub>3</sub> |
| 25          | 4443.0         | 4429           | δ<sub>sciss</sub>CH<sub>2</sub> + ν<sub>a</sub>CH<sub>2</sub> |
| 26          | 4390.5         | 4390           | δ<sub>sciss</sub>CH<sub>2</sub> + ν<sub>a</sub>CH<sub>2</sub> |
| 27          | 4329.0         | 4356           | δ<sub>sciss</sub>CH<sub>2</sub> + ν<sub>a</sub>CH<sub>2</sub> |
| 28          | 4263.8         | 4332           | δ<sub>sciss</sub>CH<sub>2</sub> + ν<sub>a</sub>CH<sub>2</sub> |
| 29          | 4174.6         | 4180           | 2[δ<sub>twist</sub>CH<sub>2</sub>, δ<sub>p</sub>COH, δwaggCH<sub>2</sub>] + δ<sub>sciss</sub>CH<sub>2</sub> |
| 30          | 4100.0         | 4107           | τCC + ν<sub>oop</sub>COH + νOH |

### Table 6. Band assignments in NIR spectra of CD<sub>2</sub>CD<sub>2</sub>OH based on GVPT2/B2PLYP-GD3BJ/def2-TZVP/CPCM calculations. Band numbering corresponds to that presented in Figure 6A.

| Peak Number | ν<sub>exp</sub> | ν<sub>calc</sub> | Assignment (Major Contribution) |
|-------------|----------------|----------------|---------------------------------|
| 1           | 7099.0         | 7126 (i)       | 2νOH |
| 2           | 6444.0         | 6495           | ν<sub>a</sub>CD<sub>2</sub> + [ν<sub>a</sub>′′′CD<sub>2</sub>, ν<sub>a</sub>′′CD<sub>2</sub>] + [ν<sub>a</sub>′′′CD<sub>3</sub>, ν<sub>a</sub>′′CD<sub>2</sub>] |
| 3           | 6232.1         | 6244           | 2δ<sub>p</sub>COH + νOH |
| 4           | 6162.9         | 6224           | 2[νCC, δwaggCD<sub>2</sub>] + νOH |
| 5           | 6063.0         | 6059           | [δ<sub>sciss</sub>CD<sub>2</sub>, νCO] + [νCC, δwaggCD<sub>2</sub>] + νOH |
| 6           | 5838.3         | 5861           | [ν<sub>a</sub>′′′CD<sub>2</sub>, ν<sub>a</sub>′′′CD<sub>3</sub>] + νOH |
| 7           | 5732.3         | 5746           | [δ<sub>rock</sub>CD<sub>2</sub>, δ<sub>rock</sub>CD<sub>3</sub>] + [δ<sub>CD3</sub>, δwaggCD<sub>2</sub>] + νOH |
| 8           | 5478.7         | 5488           | [νCC, δwaggCD<sub>2</sub>] + ν<sub>a</sub>CD<sub>3</sub> + [ν<sub>a</sub>′′CD<sub>2</sub>, ν<sub>a</sub>′′CD<sub>3</sub>] |
| 9           | 5285.8         | 5247           | ν<sub>a</sub>CD<sub>3</sub> + ν<sub>a</sub>′′CD<sub>2</sub> + νCO |
| 10          | 5160.0         | 5103           | [δ<sub>twist</sub>CD<sub>2</sub>, δ<sub>rock</sub>CD<sub>3</sub>, δ<sub>rock</sub>CD<sub>2</sub>] + 2[νCC, δwaggCD<sub>2</sub>] |
| 11          | 4903.7         | 4947           | δ<sub>p</sub>COH + νOH |
| 12          | 4797.2         | 4766           | [δ<sub>sciss</sub>CD<sub>2</sub>, νCO] + νOH |
| 13          | 4701.4         | 4714           | [δ<sub>a</sub>′′′CD<sub>3</sub>, δ<sub>a</sub>′′CD<sub>3</sub>] + νOH |
| 14          | 4588.4         | 4604           | [νCO, δwaggCD<sub>2</sub>] + νOH |
| 15          | 4525.0         | 4539           | [δ<sub>twist</sub>CD<sub>2</sub>, δ<sub>rock</sub>CD<sub>3</sub>, δ<sub>rock</sub>CD<sub>2</sub>] + δ<sub>p</sub>COH + [ν<sub>a</sub>′′′CD<sub>3</sub>, ν<sub>a</sub>′′CD<sub>2</sub>] |
| 16          | 4306.9         | 4499           | 2δ<sub>sciss</sub>CD<sub>2</sub>CO + ν<sub>a</sub>CD<sub>3</sub> |
| 17          | 4430.0         | 4447           | [ν<sub>a</sub>′′CD<sub>2</sub>, ν<sub>a</sub>′′CD<sub>3</sub>] + ν<sub>a</sub>′′CD<sub>3</sub> + 2[ν<sub>a</sub>′′CD<sub>3</sub>, ν<sub>a</sub>′′CD<sub>2</sub>] |
| 18          | 4409.6         | 4420           | [ν<sub>a</sub>′′CD<sub>2</sub>, ν<sub>a</sub>′′CD<sub>3</sub>] + [ν<sub>a</sub>′′′CD<sub>3</sub>, ν<sub>a</sub>′′CD<sub>2</sub>] |
| 19          | 4332.0         | 4334           | 2ν<sub>a</sub>′′′CD<sub>2</sub> |
| 20          | 4267.4         | 4235           | [ν<sub>a</sub>′′′CD<sub>2</sub>, ν<sub>a</sub>′′CD<sub>3</sub>] + ν<sub>a</sub>′′CD<sub>2</sub> |
| 21          | 4156.8         | 4140           | 2δ<sub>oop</sub>COH + νOH |
| 22          | 4013.0         | 4012           | δ<sub>sciss</sub>CD<sub>2</sub>CO + νOH |

### Table 7. Band assignments in NIR spectra of CD<sub>2</sub>CD<sub>2</sub>OD based on GVPT2/B2PLYP-GD3BJ/def2-TZVP/CPCM calculations. Band numbering corresponds to that presented in Figure 6B.

| Peak Number | ν<sub>Exp</sub> | ν<sub>Calc</sub> | Assignment (Major Contribution) |
|-------------|----------------|----------------|---------------------------------|
| 1           | 7771.0         | 7799           | 3νOD |
| 2           | 6468.0         | 6525           | ν<sub>a</sub>CD<sub>3</sub> + [ν<sub>a</sub>′′′CD<sub>2</sub>, ν<sub>a</sub>′′CD<sub>2</sub>] + ν<sub>a</sub>′′′CD<sub>3</sub> |
| 3           | 6450.0         | 6447           | 3ν<sub>a</sub>′′′CD<sub>2</sub> |
| 4           | 6290.1         | 6267           | 2ν<sub>a</sub>′′′CD<sub>2</sub> + ν<sub>a</sub>′′CD<sub>2</sub> |
| 5           | 5276.0         | 5289           | 2νOD |
| 6           | 4948.0         | 5020           | [νCO, δ<sub>twist</sub>CD<sub>2</sub>] + 2[δ<sub>sciss</sub>CD<sub>2</sub>, νCO] |
| 7           | 4902.2         | 4929           | [δ<sub>twist</sub>CD<sub>2</sub>, δ<sub>rock</sub>CD<sub>3</sub>] + 2[δ<sub>sciss</sub>CD<sub>2</sub>, νCO] |
| 8           | 4779.2         | 4795           | ν<sub>a</sub>′′′CD<sub>2</sub> + νOD |
Table 8. Band assignments in NIR spectra of CH₃CD₂OD based on GVPT2/B2PLYP-GD3BJ/def2-TZVP//CPMC calculations. Band numbering corresponds to that presented Figure 6C.

| Peak Number | νCalc | Assignment (Major Contribution) |
|-------------|-------|--------------------------------|
| 1           | 8781  | 3νas'CH₃                       |
| 2           | 8722  | νsCH₃ + νasCH₃ + νas'CH₃       |
| 3           | 8666  | 2νsCH₃ + νasCH₃ + νasCD₂ + νas'CH₃ |
| 4           | 7802  | 3νOD                          |
| 5           | 7257  | [δasCH₃, δas'CH₃] + νsCH₃ + νasCH₃ |
| 6           | 6445  | 3νasCD₂                       |
| 7           | 6188  | τCC + νasCH₃ + νas'CH₃        |
| 8           | 5943  | 2νasCH₃, νasCD₂ + νas'CH₃     |
| 9           | 5890  | 2[δas'CH₃, δasCH₃] + νas'CH₃ + 2νasCH₃ |
| 10          | 5845  | νsCD₂ + νasCD₂ + νasCH₃       |
| 11          | 5728  | 2νsCD₂ + νasCH₃               |
| 12          | 5288  | 2νOD                         |
| 13          | 5182  | νasCD₂ + νasCD₂ + νasCH₃     |
| 14          | 5120  | νasCD₂ + νasCH₃               |
| 15          | 5097  | νsCD₂ + νasCD₂ + νasCH₃       |
| 16          | 5046  | τCC + 2[δasCH₃, δas'CH₃]      |
| 17          | 4796  | δasCD₂ + 2[δasCH₃, δas'CH₃]   |
| 18          | 4434  | τCC + [δwaggCD₂, τCC] + νasCH₃ |
| 19          | 4372  | δsCH₃ + νasCD₂ + δsCH₃ + νas'CH₃ |
| 20          | 4341  | 2νasCD₂                       |
| 21          | 4289  | τCC + δasCD₂ + δ[δasCD₂, τCO] |
| 22          | 4233  | νsCD₂ + νasCD₂ + νasCH₃       |
| 23          | 4190  | 2νsCD₂, [δwaggCD₂, νCC] + νasCH₃ |
| 24          | 4149  | [δasCD₂, τCO] + νasCH₃, [δasCD₂, νCO] + νas'CH₃ |
| 25          | 4113  | δasCD₂ + νasCD₂ + νasCH₃     |
| 26          | 4089  | [δasCD₂, νCO] + νasCH₃        |
| 27          | 4056  | δasCD₂ + νasCH₃ + νasCH₃      |

Table 9. Band assignments in NIR spectra of CD₃CH₂OD based on GVPT2/B2PLYP-GD3BJ/def2-TZVP//CPMC calculations. Band numbering corresponds to that presented Figure 6D.

| Peak Number | νCalc | Assignment (Major Contribution) |
|-------------|-------|--------------------------------|
| 1           | 8560  | 3νasCH₂, 2νsCH₂ + νasCH₂       |
| 2           | 8508  | 3νasCH₂                       |
| 3           | 8305  | 2νsCH₂ + νasCH₂               |
| 4           | 7802  | 3νOD                         |
| 5           | 7088  | δasCD₂ + νsCD₂ + νasCH₂       |
| 6           | 6736  | δsCD₂ + νsCD₂ + νasCD₂        |
| 7           | 6706  | [νCO, δasCD₂] + νasCH₂ + νasCH₂ |
| 8           | 6614  | 3νasCD₂                       |
| 9           | 6505  | νsCD₂ + νasCD₂ + νas'CD₃      |
| 10          | 6382  | 2νsCD₂ + νasCD₂ + 2νsCD₂ + νasCD₂ |
| 11          | 5821  | [δasCD₂, τCC] + νasCH₂ + νsCH₂ |
| 12          | 5785  | 2νasCH₂, δwaggCH₂ + δasCD₂ + νasCH₂ |
Table 9. Cont.

| Peak Number | υ_{calc} | Assignment (Major Contribution) |
|-------------|----------|---------------------------------|
| 13          | 5638     | 2υ_sCH₂, υ_aCH₂ + υ_asCH₂       |
| 14          | 5547     | υ_OD + υ_sCH₂                   |
| 15          | 5473     | δ_rockCH₂ + δ_{sciss}CH₂ + υ_asCH₂ |
| 16          | 5288     | 2υ_OD                           |
| 17          | 5106     | [ν_CC, δ_{oop}CD2] + 2δ_{wagg}CH₂ |
| 18          | 4986     | 2[ν_CC, δ_sCD₂] + ν_OD           |
| 19          | 4585     | δ_{rock}CD₃ + 2δ_{wagg}CH₂       |
| 20          | 4503     | [ν_CC, δ_{oop}CD2] + δ_{wagg}CH₂ + υ_asCH₂ |
| 21          | 4459     | 2υ_asCD₃                        |
| 22          | 4431     | [δ_{as}CD₃, νCO] + [ν_CC, δ_sCD₂] + υ_as'CD₃ |
| 23          | 4385     | δ_{sciss}CH₂ + υ_asCH₂          |
| 24          | 4324     | 2υ_asCD₃, δ_{sciss}CH₂ + υ_aCH₂  |
| 25          | 4237     | δ_{wagg}CH₂ + υ_aCH₂            |
| 26          | 4167     | δ_{twist}CH₂ + υ_aCH₂           |
| 27          | 4112     | δ_{twist}CH₂ + υ_sCH₂           |
| 28          | 4015     | [δ_{rock}CD₃, δ_{rock}CH₂] + δ_{rock}CH₂ + υ_a'CD₃ |

Another insight, which becomes possible only through theoretical simulation of NIR spectra, is estimation of the relative contributions from different kinds of vibrational transitions (Table 10). As compared with methanol [41], ethanol offers better opportunity to analyze these contributions, because of higher number of isotopomers and more complex NIR spectra. The effect of various kinds of isotopic substitution of the CH₃, CH₂, and OH groups on NIR spectra may be elucidated. In the 10,000–4000 cm⁻¹ region two quanta transitions, first overtones (2υ_x) and binary combinations (υ_x + υ_y), are the most meaningful components of the spectra. In particular, binary combinations from the CH₃ group have significant contribution—e.g., for CH₃CH₂OH they are responsible for 47% of NIR intensity—while upon deuteration of the CH₃ group this contribution decreases to 32.6%. An even more pronounced effect is observed for OD derivatives, the analogous values for CH₃CH₂OD and CD₃CH₂OD are 51.2% and 35.5%, respectively. Simultaneously, the isotopic substitution of the methyl group increases the relative intensity of the first overtones, while the intensity of the second overtones remains insignificant. As expected, the importance of the second overtones increases in the upper NIR region (10,000–7500 cm⁻¹). Interestingly, this trend is not observed for the ternary combinations (υ_x + υ_y + υ_z and 2υ_x + υ_y), although for OD derivatives the 2υ_x + υ_y contribution increases and the υ_x + υ_y + υ_z contribution decreases upon deuteration of the CH₃ group. The isotopic substitution of the CH₂ group provides similar changes, but is noticeably less significant.

As can be seen (Table 10), the region above 7500 cm⁻¹ is contributed only by three and higher quanta transitions. Therefore, in this region the effect of isotopic substitution is even more visible. The deuteration of the CH₃ group increases the contributions from the second overtones at the expense of υ_x + υ_y + υ_z combinations, while the contributions from 2υ_x + υ_y remain similar. Interestingly, NIR spectrum of CD₃CD₂OD above 7500 cm⁻¹ includes the second overtones only.

These observations remain in agreement with our previous findings on methanol isotopomers [41]. However, the contributions from the three quanta transitions are more important for ethanol. For CH₃CH₂OH these transitions involve 25.9% of total intensity (10,000–4000 cm⁻¹), while for CH₃OH this value was found to be 19.2%. The difference between CD₃OD and CD₃CD₂OD is even larger (23.5% vs. 36.7%).
Table 10. Contributions (in %) from the first and second overtones as well as binary and ternary combinations into NIR spectra of ethanol isotopomers based on GVPT2//B2PLYP-GD3BJ//def2-TZVP//CPCM calculations. \(^{a}\)

|                | 10,000–4000 cm\(^{-1}\) | 10,000–7500 cm\(^{-1}\) | 7500–4000 cm\(^{-1}\) |
|----------------|--------------------------|---------------------------|------------------------|
|                | 2\(v_x\) | 3\(v_x\) | \(v_x + v_y\) | \(v_x + v_y + v_z\) | 2\(v_x + v_y\) | 2\(v_x\) | 3\(v_x\) | \(v_x + v_y\) | \(v_x + v_y + v_z\) | 2\(v_x + v_y\) | 2\(v_x\) | 3\(v_x\) | \(v_x + v_y\) | \(v_x + v_y + v_z\) | 2\(v_x + v_y\) |
| CH\(_3\)CH\(_2\)OH | 26.1     | 1.7     | 47.0     | 14.3     | 10.9     | 0.0     | 39.7     | 0.0     | 22.3     | 38.0     |
| CH\(_3\)CH\(_2\)OD | 18.0     | 2.2     | 51.2     | 17.4     | 11.1     | 0.0     | 55.5     | 0.0     | 15.4     | 29.1     |
| CH\(_3\)CD\(_2\)OH | 35.8     | 1.7     | 41.5     | 11.8     | 9.2      | 0.0     | 43.9     | 0.0     | 30.5     | 25.6     |
| CD\(_3\)CH\(_2\)OH | 40.9     | 1.2     | 32.6     | 15.1     | 10.1     | 0.0     | 66.9     | 0.0     | 1.4      | 31.7     |
| CD\(_3\)CD\(_2\)OH | 46.0     | 0.3     | 23.7     | 15.8     | 14.2     | 0.0     | 0.0      | 0.0     | 0.0      | 0.0      |
| CD\(_3\)CD\(_2\)OD | 43.1     | 2.0     | 19.2     | 17.8     | 17.9     | 0.0     | 69.9     | 0.0     | 16.7     | 13.4     |
| CD\(_3\)CD\(_2\)OH | 27.9     | 3.2     | 44.7     | 15.0     | 9.2      | 0.0     | 76.3     | 0.0     | 0.5      | 23.2     |

\(^{a}\) The comparison is based on integrated intensity (cm\(^{-1}\)) summed over simulated bands, convoluted with the use of Cauchy–Gauss product function (details in the text) in relation to the total integrated intensity.

3. Experimental and Computational Methods

3.1. Materials and Spectroscopic Measurements

In Table 11 are collected the details on the samples used in this work. The experimental spectrum of CH\(_3\)CH\(_2\)OH was taken from our previous work \([29]\). All samples were used as received, while solvent (CCl\(_4\)) was distilled and additionally dried using freshly activated molecular sieves (Aldrich, 4A). All ethanols were measured in CCl\(_4\) solution (0.1 mol dm\(^{-3}\)). NIR spectra were recorded on Thermo Scientific Nicolet iS50 spectrometer using InGaAs detector, with a resolution of 2 cm\(^{-1}\) (128 scans), in a quartz cells (Hellma QX, Hellma Optik GmbH, Jena, Germany) of 100 mm thicknesses at 298 K (25 °C).
Table 11. Samples used in this study

| Sample         | Purity | D Atom Content | Other Remarks |
|----------------|--------|----------------|---------------|
| 1 CH₃CH₂OD     | 99%    | ≥99.5%         |               |
| 2 CH₃CD₂OH     | 99%    | 98%            |               |
| 3 CD₃CH₂OH     | 99%    | 99%            |               |
| 4 CD₃CD₂OH     | 99%    | 99.5%          |               |
| 5 CD₃CD₂OD     | >99%   | ≥99.5%         | anhydrous     |
| 6 CCl₄         | >99%   | -              |               |

Samples were purchased from Sigma-Aldrich Chemie GmbH (Taufkirchen, Germany).

3.2. Computational Procedures

Our calculations were based on density functional theory (DFT) with double-hybrid B2PLYP density functional [55] (unfrozen core) coupled with Karlsruhe triple-ζ valence with polarization (def2-TZVP) [56] basis set. Grimme’s third formulation of empirical correction for dispersion with Becke-Johnson damping (GD3BJ) was applied [57]. To better reflect solvation of molecules, CCl₄ cavity in solvent reaction field (SCRF) [58] was included at conductor-like polarizable continuum (CPCM) [59] level. Very tight criteria for geometry optimization and 10⁻¹⁰ convergence criterion in SCF procedure were set. Electron integrals and solving coupled perturbed Hartree-Fock (CPHF) equations were calculated over a superfine grid. The selected method provided good reproduction of NIR spectra of various molecules in CCl₄ solution [19,29,30].

We carried out the anharmonic vibrational analysis at generalized vibrational second-order perturbation theory (GVPT2) [60,61] level. In this approach, the anharmonic frequencies and intensities of the vibrational transitions up to three quanta were obtained. This allows to simulate fundamental, first and second overtones, as well as binary and ternary combination bands. Quantum mechanical calculations were carried out with Gaussian 16 (A.03) [62]. One of the major features implemented in GVPT2 approach is the automatic treatment of tight vibrational degenerations, i.e., resonances [63]. In this work the search for resonances included Fermi (i.e., 1-2) of type I (ω₁ ≈ 2ω₂) and type II (ω₁ ≈ ω₂ + ω₃), and Darling–Dennison (i.e., 2-2, 1-1, and 1-3) resonances. All possible resonant terms within search thresholds were included in the variational treatment. The resonance search thresholds (respectively, maximum frequency difference and minimum difference PT2 vs. variational treatment; in (cm⁻¹)) were: 200 and 1 (for the search of 1–2 resonances), 100 and 10 (for 2-2, 1-1, and 1-3).

To display the simulated spectra we applied a four-parameter Cauchy–Gauss (Lorentz–Gauss) product function [20]. The theoretical bands were modelled with $a_2$ and $a_4$ parameters equal to 0.055 and 0.015, resulting with full-width at half-height (FWHH) of 25 cm⁻¹. Exception was made for better agreement with the weaker and broader experimental bands, which are presented in Figures 3 and 4. In this case the values were 0.075, 0.015, and 35 cm⁻¹, respectively. The final theoretical spectra were obtained by combining the spectra of trans and gauche conformers, mixed in accordance with the calculated abundances of each form [64]. The relative abundances of the gauche ($n_g$) and trans ($n_t$) conformers were determined as following equation [65].

$$\frac{n_g}{n_t} = \frac{A_t}{A_g} e^{-\Delta G^{298}/RT}$$

where Gibbs free energy ($\Delta G$) corresponds to the value calculated at 298 K corrected by anharmonic (VPT2) zero-point energy (ZPE); $A_t$ and $A_g$ are the degeneracy prefactors of the Boltzmann term for the gauche (1) and trans (2) conformers.

The band assignments were aided by calculations of potential energy distributions (PEDs). PEDs were obtained with Gar2Ped software [66], using natural internal coordinate system defined in accordance with Pulay [67]. The numerical analysis of the theoretical results and the processing of the experimental spectra were performed with MATLAB R2016b (The Math Works Inc.) [68].
4. Conclusions

Isotopic substitution leads to much higher variability in NIR spectra as compared with IR spectra, due to significant contribution from the combination bands. The pattern of OH/OD, CH$_3$/CD$_3$, and CH$_2$/CD$_2$ groups in ethanol often leads to fine spectral changes, which may be monitored and explained in detail by anharmonic quantum mechanical simulations. Our studies were devoted to NIR spectra of eight isotopomers of ethanol (C$_X$$_3$C$_X$$_2$O$_X$ (X = H, D)) by using anharmonic GVPT2 vibrational analysis. The calculations were performed at several levels of electronic theory, including DFT and MP2 to find accurate and efficient theoretical approach for studies of isotopic effect in NIR spectra. Our results indicate that DFT approach using double-hybrid B2PLYP functional, coupled with def2-TZVP basis set, and supported by GD3BJ correction with CPCM solvent model yielded the best results. The theoretical spectra obtained by this approach enabled us to assign most of NIR bands, including two ($2\nu_x$ and $\nu_x + \nu_y$) and three quanta ($3\nu_x$, $\nu_x + \nu_y + \nu_z$, and $2\nu_x + \nu_y$) transitions. Accuracy of these calculations permitted us to analyze theoretical NIR spectra of CH$_3$CD$_2$OD and CD$_3$CH$_2$OD for which the experimental spectra are not available. The effect of the isotopic substitution of the OH, CH$_3$, and CH$_2$ groups was satisfactory reproduced and explained. Moreover, the relative contributions of selected groups and kinds of transitions were elucidated and discussed. The contributions from the CH$_3$ group appear to be more important than those from the CH$_2$ group. The isotopic substitution in the CH$_3$ group leads to the most prominent intensity changes in NIR spectra as compared to the changes due to the substitution of the other groups. The bands from the three quanta transitions are more important for isotopomers of ethanol than for derivatives of methanol.

Supplementary Materials: The following are available online, Figures S1–S35; Tables S1–S9.

Author Contributions: Conceptualization, K.B.B. and M.A.C.; Methodology K.B.B. and J.G.; Formal analysis, J.G. and K.B.B.; Investigation, all authors; Writing—original draft preparation, K.B.B.; Writing—review and editing, all authors; Supervision, C.W.H. and M.A.C.

Funding: This work was supported by the National Science Center Poland, Grant No. 2017/27/B/ST4/00948.

Acknowledgments: Calculations have been carried out in Wroclaw Centre for Networking and Supercomputing (http://www.wcss.pl), under grant no. 163.

Conflicts of Interest: The authors declare no conflict of interest. The founders had no role in the design of the study; in the collection, analyses, or interpretation of data; in the writing of the manuscript, or in the decision to publish the results.

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Sample Availability: Samples of the compounds (CH$_3$CH$_2$OD, CH$_3$CD$_2$OH, CD$_3$CH$_2$OH, CD$_3$CD$_2$OH, CD$_2$CD$_2$OD) are available from the authors.

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