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Citation
Deutscher, N. M., D. W. T. Griffith, G. W. Bryant, P. O. Wennberg, G. C. Toon, R. A. Washenfelder, G. Keppel-Aleks, et al. 2010. Total Column \((\text{CO}_2)\) Measurements at Darwin, Australia – Site Description and Calibration Against In Situ Aircraft Profiles. Atmospheric Measurement Techniques 3, no. 4: 947–958.

Published Version
doi:10.5194/amt-3-947-2010

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Accessibility
Total column CO\textsubscript{2} measurements at Darwin, Australia – site description and calibration against in situ aircraft profiles

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Received: 3 March 2010 – Published in Atmos. Meas. Tech. Discuss.: 17 March 2010
Revised: 26 May 2010 – Accepted: 7 July 2010 – Published: 19 July 2010

Abstract. An automated Fourier Transform Spectroscopic (FTS) solar observatory was established in Darwin, Australia in August 2005. The laboratory is part of the Total Carbon Column Observing Network, and measures atmospheric column abundances of CO\textsubscript{2} and O\textsubscript{2} and other gases. Measured CO\textsubscript{2} columns were calibrated against integrated aircraft profiles obtained during the TWP-ICE campaign in January–February 2006, and show good agreement with calibrations for a similar instrument in Park Falls, Wisconsin. A clear-sky low airmass relative precision of 0.1% is demonstrated in the CO\textsubscript{2} and O\textsubscript{2} retrieved column-averaged volume mixing ratios. The 1% negative bias in the FTS X\textsubscript{CO\textsubscript{2}} relative to the World Meteorological Organization (WMO) calibrated in situ scale is within the uncertainties of the NIR spectroscopy and analysis.

1 Introduction

Carbon dioxide in the atmosphere is the most important contributor to positive radiative forcing responsible for the enhanced greenhouse effect (Forster et al., 2007). To better understand and manage CO\textsubscript{2} emissions, estimates of source and sink strength and variability are required on at least regional scales. In the past, these regional strengths and patterns of surface exchange across the globe have been quantified via a combination of in situ CO\textsubscript{2} measurements from a global network of surface sites (GLOBALVIEW-CO\textsubscript{2}), and inverse modeling studies (Gurney et al., 2002). Despite the high precision and accuracy of the in situ measurements, this approach exhibits a high variability of source/sink distributions across different transport models (Gurney et al., 2003), particularly when attempting to constrain regions on a small spatial scale. Errors in modeling transport are a contributing factor to the varying results between models; however errors are also introduced due to rectifier effects, which arise because the surface fluxes are seasonally and diurnally correlated with vertical transport (Denning et al., 1996; Gurney et al., 2002). Higher surface concentrations occur at night-time and in winter due to CO\textsubscript{2} build up in a shallow planetary
vertical-transport-induced variability compared to surface data, though at the disadvantage of being less sensitive than in situ measurements to surface sources and sinks. Column measurements can therefore complement the existing in situ network; however, high precision in the column measurements is required for low variability sites, such as those in the Southern Hemisphere, in order to determine meaningful seasonal and spatial patterns (Rayner and O’Brien, 2001). Simulations have confirmed their potential for constraining the global carbon balance if obtained with suitable precision and accuracy (e.g., precision better than 2.5 µmol mol⁻¹ for monthly average column data with global coverage with a 8° × 10° footprint) (Rayner and O’Brien, 2001). Particular care must be taken, however, with systematic biases, which must be limited to a few tenths of a µmol mol⁻¹ (Chevallier et al., 2007).

The Total Carbon Column Observing Network (TCCON) (http://www.tccon.caltech.edu; Wunch et al., 2010a) is a network of ground-based solar Fourier Transform Spectrometers (FTS) operating in the near infrared (NIR) spectral region, focused on providing highly accurate and precise column measurements of CO₂, as well as simultaneously measuring O₂, CH₄, H₂O, HDO, HF, CO and N₂O and other species. The column data resulting from the site considered here and sites with similar setups can be used not only for validation of satellites, but also in atmospheric inversion modeling studies. Quantification of biases and errors is extremely important for inversion studies, so calibration and subsequent quantification of precision and accuracy of the column measurements is important for their use in this application. Column measurements are especially important in the tropics, as convection is consistently strong, and as a result flux signals are only weakly seen in surface measurements (Gloor et al., 2000; Rayner and O’Brien, 2001). The importance of tropical column measurements in constraining regional flux estimates has previously been emphasized in modeling studies (Gloor et al., 2000; Olsen and Randerson, 2004; Rayner and O’Brien, 2001). The TCCON can potentially provide these data.

By measuring the same quantities (i.e., column abundances) in the same spectral regions as satellite-borne instruments such as the Orbiting Carbon Observatory (OCO) (Crisp et al., 2004), the Greenhouse Gases Observing Satellite (GOSAT) (Inoue et al., 2006) and the Scanning Imaging Absorption SpectroMeter for Atmospheric CHartography (SCIAMACHY) (Bovensmann et al., 1999), TCCON will be able to provide a definitive validation and calibration of the satellite products. The global network of sites comprising TCCON will be particularly useful for determining and calibrating spatial and temporal variation in the satellite products. A global network is necessary to provide satellite validation over a range of latitudes, surface types, environments and airmasses.

This paper describes the dedicated solar observatory deployed to Darwin, Australia (12.425°S, 130.891°E) in August 2005. The observatory currently is the only continuously operating solar FTS located within the tropics, making it crucial for providing data to constrain this important region, as well as to provide calibration and validation of satellite data over the tropics. Constraining regional-scale source and sink strengths in the tropics is difficult because there are limited CO₂ measurements in this region. Quantifying tropical CO₂ sources and sinks is important because of carbon stocks in tropical rainforests, deforestation, tropical wetlands, the unique presence of savannahs, and the large contribution from regular biomass burning. Darwin has a very distinct climatic pattern—a monsoonal wet season from December to March, a dry season from May to September, and transitions in between these times.

In this paper Sect. 2 describe the instrumentation, Sect. 3 the measurement site and Sect. 4 introduces the data analysis. In Sect. 5 we outline the calculation of total column mole fraction and investigate the reproducibility of these measurements in Sect. 6. Section 7 presents the calibration of atmospheric CO₂ columns to the global reference scale via comparison to integrated in situ aircraft profiles obtained during the Tropical Warm Pool – International Cloud Experiment (TWP-ICE) in January – February 2006. Section 8 contains a discussion of the errors. This is followed by the summary and conclusions, and Appendix A contains a method to correct for airmass dependent effects.

2 Instrumentation

The solar observatory was assembled at the California Institute of Technology (Caltech), and shipped to Darwin, Australia in 2005. The observatory is very similar to that detailed by Washenfelder et al. (2006), who describe the first dedicated TCCON solar observatory deployed to Park Falls, Wisconsin. In summary, a Bruker IFS 125/HR FTS (Bruker Optics, Germany) is mounted inside a customized 6.1×2.4×2.6 m shipping container. The spectrometer has two room-temperature NIR detectors, an indium gallium arsenide (InGaAs – 3800 – 12 000 cm⁻¹) and a silicon diode (Si – 9500 – 30 000 cm⁻¹) which operate in simultaneous dual acquisition mode with a 10,000 cm⁻¹ dichroic filter and calcium fluoride (CaF₂) beamsplitter. A red filter limits the Si spectra to 15 600 cm⁻¹ to avoid aliasing. Atmospheric column abundances of CO₂, O₂, CH₄, H₂O, HDO, HF, CO and N₂O can be retrieved from the NIR spectra. A Bruker Optics solar tracker is mounted on the roof of the container and
protected by an astronomy dome (ROBODOME™, Technical Innovations Inc., Gaithersburg, Maryland). The tracker directs the solar beam through a CaF$_2$ window in the roof of the container and onto a folding mirror below, where it is then reflected into the spectrometer. A pickoff mirror directs a small fraction of the beam to a Si quadrant detector that controls the active solar tracking. Scans are routinely collected over optical path differences of −5 cm to 45 cm, with a mirror speed of 0.633 cm s$^{-1}$ (fringe rate 10 kHz, sample rate 20 kHz) resulting in a collection time of 76 s for a single scan.

Daily operation of the container is fully automated. A weather station monitoring temperature, relative humidity, solar radiation, presence of rain, leaf wetness and pressure is also mounted on the roof of the container. Housekeeping data, including meteorological parameters, are logged continuously at a frequency of 1 Hz. Active solar tracking and spectral collection are initiated when the solar intensity measured at the quadrant detector is greater than ∼45 W m$^{-2}$. The pressure as measured with a Setra Systems Inc. Model 270 pressure transducer is periodically calibrated against a mercury manometer mounted inside the container. The precision of the pressure transducer is approximately 0.03% (0.3 hPa).

The differences from the Park Falls setup detailed in Washenfelder et al. (2006) are the inclusion of a leaf wetness sensor on the weather station (now also installed at Park Falls) and the increased fringe rate from 7.5 kHz to 10 kHz (which has since been changed back to 7.5 kHz to reduce data volume) and hence reduced spectral collection time of 76 s rather than 105 s. On 29 September 2005, the method of data collection was changed from the normal AC recording to DC acquisition mode. Collecting DC interferograms enables correction for source brightness fluctuations during each scan that occur due to changing thin cloud and aerosol cover. The switch to DC acquisition improves the reproducibility of column-average dry-air mole fraction (X$_{CO_2}$) under partly cloudy conditions, as well as the number of usable spectra obtained (Keppel-Aleks et al., 2007).

In addition to the room-temperature InGaAs and Si-diode detectors, a liquid nitrogen cooled indium antimonide (InSb) detector (2000–6500 cm$^{-1}$) is installed. Spectra are collected with this detector at times of biomass burning that occur during the dry season in the savannahs across tropical Australia. The additional spectral coverage provided by the InSb detector results in the ability to simultaneously monitor biomass burning species including CO, HCN, C$_2$H$_6$, C$_3$H$_2$ and H$_2$CO (Paton-Walsh et al., 2010), although these gases are not the subject of this paper.

The instrument has operated almost continuously since installation, however short data gaps of up to one month have occurred due to a battery failure in April 2006 and computer failure in June 2007. Aside from this, the major problem encountered has been degradation of the external gold-coated solar tracker mirrors, which have been replaced on two occasions, in early December 2005 and early November 2006. The mirror coating degradation causes a progressive reduction in light throughput, resulting in decreased signal-to-noise ratio, and consequently, measurement reproducibility. In the first two years of operation, data were collected on 588 days, with an average of 250 spectra per day during the dry season.

3 Measurement site

The laboratory is located at the Department of Energy (DoE) Tropical Western Pacific (TWP) Atmospheric Radiation Measurement (ARM) site, located adjacent to the Australian Bureau of Meteorology (BoM) station at Darwin International Airport (12.425 S, 130.89 E, 30 m above sea level). The site is 9 km from the city of Darwin (population 90 000). The airport typically has fewer than 30 incoming and outgoing commercial flights per day plus sporadic military use. Thus interference from aircraft exhaust emissions is minimal, especially as airport peak traffic is after sunset, when solar measurements cannot be collected. The ARM site was chosen because of the existing infrastructure and technical support available, as well as the suite of ancillary measurements available both there and at the BoM site. An in situ FTIR gas analyzer similar to that described by Esler et al. (2000a, b) has been in place since February 2007, measuring CO$_2$, $\delta^{13}$C in CO$_2$, CH$_4$, N$_2$O and CO at 12 m above ground level.

4 Data analysis

DC corrected spectra are analyzed using GFIT, a non-linear least-squares fitting algorithm developed at the Jet Propulsion Laboratory (G. Toon). A theoretical atmospheric transmittance spectrum is calculated using molecular absorption coefficients, calculated atmospheric ray paths, model assimilated profiles of temperature, pressure and humidity and assumed gas dry-air mole fraction profiles. The temperature, pressure and water vapor profiles are obtained from National Center for Environmental Prediction (NCEP) reanalysis data (Kalnay et al., 1996) provided by the NOAA/ESRL Physical Sciences Division, and interpolated in time and space from six hourly data to local solar noon and site latitude/longitude. The a priori CO$_2$ profile is based on a model fitted to GLOBALVIEW data. The GFIT forward model calculation uses 70 vertical levels spaced at 1 km intervals to represent the atmosphere. The calculated and measured spectra are compared, and the RMS difference between the two is minimized by iteratively scaling the gas VMR profiles.

CO$_2$ columns are retrieved in two bands centered at 6228 cm$^{-1}$ ($v_1 + 4v_2 + v_3$) and 6348 cm$^{-1}$ ($2v_1 + v_2 + 2v_3$), from which we use the average column amount. Figure 1a shows a typical pair of simultaneously acquired InGaAs and Si spectra, with (b) an expanded view of the 6228 cm$^{-1}$
Fig. 1. (a) Typical InGaAs (black) and Si (grey) spectra, obtained simultaneously from Darwin on July 10, 2006. (b) Expanded view of the CO$_2$ 6228 cm$^{-1}$ ($\nu_1 + 4\nu_2 + \nu_3$) and 6348 cm$^{-1}$ ($2\nu_1 + \nu_2 + 2\nu_3$) bands from the same spectrum.

and 6348 cm$^{-1}$ CO$_2$ bands. The O$_2$ A-band (13070 cm$^{-1}$), which was to be used for retrievals by the OCO satellite (Crisp et al., 2004), is not used in our analyses. Instead, we use the 7882 cm$^{-1}$ O$_2$ band, which has the advantage of being collected using the same detector as the two CO$_2$ bands, and being closer in frequency, so that systematic effects such as instrument lineshape (ILS) errors, zero level offset, and aerosol and cloud scattering will be more similar between the three bands and therefore mostly cancel in the CO$_2$/O$_2$ ratio.

The spectral line parameters are based on the HITRAN linelist (Rothman et al., 2005), with improvements to the O$_2$ 7882 cm$^{-1}$ band (Newman et al., 2000; Yang et al., 2005), CO$_2$ 6228 cm$^{-1}$ and 6348 cm$^{-1}$ bands (Toth et al., 2006, 2010), and being closer in frequency, so that systematic effects such as instrument lineshape (ILS) errors, zero level offset, and aerosol and cloud scattering will be more similar between the three bands and therefore mostly cancel in the CO$_2$/O$_2$ ratio.

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5 Calculation of total column mixing ratios

We convert the raw retrieved whole air columns to column-average dry-air mole fractions (DMFs) via Eq. (1) by dividing the vertical column of gas by that of the total dry atmosphere:

$$X_i = \frac{VC_i}{m_{air}^{dry} \tilde{g}} - VC_{H_2O} \frac{m_{H_2O}^{dry}}{m_{air}^{dry}}$$

where:

- $X_i$ = column-average dry-air mole fraction of gas $i$
- $VC_i$ = vertical column of gas $i$ (molecules cm$^{-2}$)
- $P_s$ = surface pressure (mb)
- $\tilde{g}$ = absorber weighted gravitational acceleration
- $m_{air}^{dry}$ = molecular mass of dry air (28.964 g mol$^{-1}$)
- $m_{H_2O}^{dry}$ = molecular mass of H$_2$O (18.02 g mol$^{-1}$)

For O$_2$, we know the atmospheric profile to have a constant dry-air mole fraction of 0.2095. We can therefore rearrange Eq. (1) to yield the relationship (2), which presents the total...
dry air column in terms of the vertical column of \( O_2 \), rather than pressure. Substituting back into (1) gives Eq. (3), which allows \( X_i \) to be calculated using \( V \text{CO}_2 \) from simultaneously acquired spectra. This eliminates systematic errors that are common to species \( i \) and \( O_2 \), such as those from solar intensity variation and ray path variation due to scattering and solar pointing. We also eliminate site-dependent biases that may arise in the surface pressure and gravity, thereby increasing the measurement accuracy. Systematic errors in \( X_i \) may be introduced in the long-term by any variation in retrieved \( O_2 \) that does not similarly affect trace gas \( i \).

\[
\frac{P_s}{m_{\text{air}} \cdot g} = \frac{V \text{CH}_2 \text{O}}{m_{\text{H}_2 \text{O}}} = \frac{V \text{CO}_2}{X_2} = \frac{V \text{CO}_2}{0.2095} \tag{2}
\]

\[
X_i = \frac{V \text{C}_i}{V \text{CO}_2} \cdot 0.2095 \tag{3}
\]

One of the advantages of Darwin over Park Falls is that the natural variations of \( X \text{CO}_2 \) are much smaller, allowing systematic errors in the data to be more easily identified and quantified. The high range of solar zenith angles and large number of sunny days further helps with this error characterization.

Plots of \( X \text{CO}_2 \) as a function of solar zenith angle reveal \( \sim 1\% \) larger values around noon than sunrise/set. This occurs at all TCCON sites and seasons, even when the true variation of \( \text{CO}_2 \) is known to be small. This so-called airmass-dependent-artifact is believed to arise from spectroscopic deficiencies and has been corrected in the data presented here. A description of this correction is given in Appendix A, and further detail given by Wunch et al. (2010a).

6 \( \text{O}_2 \) and \( \text{CO}_2 \) reproducibility

\( X \text{CO}_2 \) and \( X \text{O}_2 \) are filtered for relative solar intensity variation of less than ten percent during each scan, and retrieval error of less than 0.02 (\( \sim 2\% \)) in the calculated profile scale factor. \( X \text{O}_2 \) allows us to ensure the repeatability of the measurements that could otherwise be limited by scattering, ray path variation or the precision of the pressure transducer. For each day of measurements, the percentage standard deviation is calculated to give an indication of the within-day measurement reproducibility (GAW, 2007). Figure 3 shows the retrieved \( X \text{CO}_2 \) and \( X \text{O}_2 \), along with the daily mean \( \pm \) standard deviation (crosses and error bars) for two example days, June 27 and 28, 2006. Here the \( X \text{O}_2 \) is calculated via Eq. (1) and scaled to a mean value of 0.2095, but for the remainder of this manuscript the unscaled \( X \text{O}_2 \) is used. On clear sky days such as these (\( \sim 60 \) spectra), the reproducibility of \( X \text{CO}_2 \) is better than 0.09% (one standard deviation) and \( X \text{O}_2 \) better than 0.11%. The better reproducibility in \( X \text{CO}_2 \) over \( X \text{O}_2 \) occurs because of the cancellation of errors common to \( \text{CO}_2 \) and \( \text{O}_2 \) columns. The clear sky reproducibility of the column measurements is therefore of the order of 0.1%.

The reproducibility of the \( \text{CO}_2 \) and \( \text{O}_2 \) retrievals is much better than the RMS uncertainty achieved by the spectral fit, such as that illustrated in Fig. 2. Since the fitted \( \text{CO}_2 \) windows contain around 100 spectral lines it is possible to achieve a reproducibility in the retrieved columns that is better than the RMS spectral fit. The RMS residual is also affected by systematic effects which are invariant from spectrum-to-spectrum and do not affect the reproducibility of the measurement.
Fig. 5. Trajectory of the 4 February DoE-Proteus flight. The black dot represents the ARM site at which the FTS is located. The color scale shows the progression through time in UT, from blue at take off, through green at the upward spiral completion to red at landing.

A time series of 14 months of $X_{CO_2}$ data is presented in Fig. 4. A seasonal pattern in $CO_2$ is evident, with peak mole fraction at the start of the monsoon season, followed by rapid uptake due to the onset of plant growth and then a resumption of $CO_2$ growth from the end of the rainy season. Also shown in this figure is the time series of retrieved continuum level, hence indicating the relative signal levels. A decrease in signal can be seen throughout the time series. Some of the lost signal is regained with replacement solar tracker mirrors in December 2005 and November 2006.

7 Column $CO_2$ calibration against in situ measurements

We have shown that good reproducibility can be obtained on the solar column $CO_2$ measurements. In addition to the airmass-dependent biases discussed earlier, there are also systematic biases that are airmass-independent, probably resulting from spectroscopic errors in the integrated band intensities. These are actually easier to deal with because they are invariant. In order to correct for these biases and place these measurements on the same $CO_2$ scale as used for in situ measurements (Tans et al., 2003), and thus to be able to use them alongside in situ data for satellite validation and inverse modeling, they need to be calibrated against a standard technique on an absolute scale. This requires simultaneous, co-located measurement of $CO_2$ profiles that are calibrated on the WMO $CO_2$ scale.

The Tropical Warm Pool International Cloud Experiment aircraft mission (TWP-ICE, January–February 2006) provided the opportunity to perform such calibration against $CO_2$ profiles measured in situ with the “ER-2” NDIR $CO_2$ analyzer on board the Proteus aircraft. Simultaneous, co-located $CO_2$ profiles obtained during this campaign are integrated for comparison with the retrieved FTS column $CO_2$.

The Harvard University “ER-2” $CO_2$ instrument is described in detail elsewhere (Daube et al., 2002). Briefly, the instrument is based on a modified NDIR $CO_2$ analyzer (Li-Cor LI-6251), which measures the light absorption of the sample at 4.26 µm relative to a reference gas of known $CO_2$ concentration. A Nafion drier and a dry ice trap were added to the instrument for its deployment on Proteus. Dilution by water vapor and wall effects are negligible at the resulting sample dewpoint ($<−60^\circ C$). Calibrations with reference (zeroing), low- and high-span, and a long-term standard, were performed in flight every 10, 30 and 120 min, respectively. The calibration with a long-term standard ensures flight-to-flight accuracy and precision, which is augmented by ground-based calibrations of the source and instrument tanks before and after the mission. All the calibration gases are traceable to WMO primary standards. The demonstrated (>160 flights) precision and accuracy of the instrument is better than $±0.1 \mu mol \cdot mol^{-1}$. The $CO_2$ mixing ratio is reported at 0.5 Hz (2 s response time).

During TWP-ICE, $CO_2$ was measured on Proteus during five flights: 25, 27 and 29 January and 2 and 4 February 2006. These flights were all during the wet or monsoon season. Only the last flight occurred under clear sky conditions, so the comparison uses data from the Feb 4 flight only. The flight consisted of an upward and a downward spiral centered on the ARM site. $CO_2$ measurements commenced at
00:51 UT (10:21 LT). Figure 5 shows the trajectory of the 4 February flight centered on the observatory. The color scale from red to blue represents the progression through time of the flight.

The total CO₂ vertical column from the in situ aircraft profile is determined by integrating the measured profiles:

\[ V_{\text{CO}_2} = \int_0^P f_{\text{CO}_2}^\text{dry} \cdot (1 - f_{\text{H}_2\text{O}}) \cdot g(p) \cdot m \, dp \]  

(4)

with \( m = \left[ m_{\text{H}_2\text{O}} \cdot f_{\text{H}_2\text{O}} + m_{\text{air}}^\text{dry} \cdot (1 - f_{\text{H}_2\text{O}}) \right] \)  

(5)

where:

\( f_{\text{CO}_2}^\text{dry} (p) = \text{dry} – \text{air mole fraction of CO}_2 (\text{mol mol}^{-1}) \);  
\( f_{\text{H}_2\text{O}} (p) = \text{mole fraction of H}_2\text{O} (\text{mol mol}^{-1}) \);  
\( m(p) = \text{mean molecular mass of (wet) air} \);  
\( g(p) = \text{gravitational acceleration} \).

This formulation requires measurements of CO₂ and H₂O throughout the vertical profile. The Proteus Harvard instrument provides \( f_{\text{CO}_2}^\text{dry} \), but there was no reliable co-located measurement of water vapor aboard the Proteus flights. As an alternative, we use humidity data collected using a Vaisala RS92 sensor on board a balloon-borne sonde launched from the ARM site that flew between 0845 and 1030 local time on the flight day. These humidity data are corrected using the algorithm presented by Vömel et al. (2007a, b), with a further correction to eliminate solar zenith angle dependence (Hume, 2007). They are then interpolated onto the aircraft pressure grid, and used as \( f_{\text{H}_2\text{O}} \) in the integration. The contribution of H₂O to the total error budget is small – a systematic error of 5% (the quoted total uncertainty in the Vaisala RS92 for a sounding) in the balloon-borne H₂O concentrations would result in an \( X_{\text{CO}_2} \) error of 0.03 µmol mol⁻¹ (0.008%).

The 4 February airborne spiral (139–961 hPa) covers a large vertical portion of the troposphere. Nevertheless, in order to determine the entire integrated column, it is necessary to extrapolate the CO₂ profile to the lowermost part of the planetary boundary layer (PBL) and to the uppermost part of the troposphere and the stratosphere. This requires formulating some assumptions about the CO₂ profile within these regions of the atmosphere.

The average of five NCEP model runs commencing at different times is used to determine the PBL height for Darwin. The NCEP model runs agree well with the PBL height calculated from the balloon-borne soundings. It is assumed that the PBL is sufficiently well mixed during the time of the flight, so that a constant CO₂ profile is used based on aircraft PBL measurements, and scaled to match the bottom of the profile. In the PBL, \( f_{\text{CO}_2}^\text{dry} \) is particularly sensitive to spatial variations, due to localized sources/sinks. At these altitudes, the Proteus digressed from the spiral profile to make its landing. The inset in Fig. 6 shows the suite of CO₂ measurements in the PBL from Proteus during the ascent and descent, along with the modeled PBL height. The \( f_{\text{CO}_2}^\text{dry} \) value varies by approximately ± 1 µmol mol⁻¹, within the error estimate given.

The variability is both spatial and temporal, and occurs during the spirals and the preceding and succeeding flight time.

The tropopause height is supplied by NCEP reanalysis. Below the tropopause but above the aircraft ceiling, the free troposphere CO₂ profile is assumed to be equal to the mean profile value measured above 200 hPa. Above the tropopause, a stratospheric model is used to generate the stratospheric CO₂ profile. The model uses time and latitude (and N₂O as internal variable) to determine the age of the air since entering the stratosphere (Andrews et al., 1999; 2001a, b) along with surface measurements of CO₂ at Samoa and Mauna Loa (Boering et al., 1996) to compute stratospheric CO₂ profiles for either the equatorial region (~ −5° to +5°) or mid-latitudes (35° to 55°). To obtain the value at Darwin, a linear interpolation is performed between the equatorial profile at 5° latitude and the profile at 35° latitude to obtain a stratospheric CO₂ profile at 12.425° latitude. The rapid transport of tropical CO₂ to the lower Tropical Tropopause Layer (TTL) has recently been confirmed using CO₂ measurements on the WB-57 during CR-AVE (2006) and on Proteus during TWP-ICE (Park et al., 2007).

Profile values from the Proteus are allocated errors of 0.1 µmol mol⁻¹ (1-sigma) in calculating the integrated column error, while all the assumed profile values are assigned error bars of 2 µmol mol⁻¹. The measured (crosses) and extrapolated (lines) portions of the profiles are shown in Fig. 6. The entire profile is weighted by the mean vertically resolved FTS averaging kernel over the period of the overpass and both retrieval windows, before comparing with FTS column retrievals. The method of applying and calculating the averaging kernel is described by Wunch et al. (2010b). The averaging kernel is derived by perturbing the mole fraction...
at each particular level and observing the resulting change in the retrieval. In a profile scaling retrieval, the rows of the averaging kernel matrix all have the shape of the a priori profile, and therefore the NxN matrix can be simplified to a N-vector, like that illustrated in Washenfelder et al. (2006).

Table 1 summarises the integrated aircraft and FTS average $X_{CO_2}$ during the upward and downward profiles. Unlike Washenfelder et al. (2006) we use the average $CO_2$ from the two bands, and do not previously scale the $O_2$ to a mean value of 0.2095. Rather than deriving separate $CO_2$ and $O_2$ corrections, one correction factor is used to cover the whole $X_{CO_2}$ calculation. This results in identical calibrated $X_{CO_2}$ numbers, but simplifies the calibration process. Ratiosing the FTS and aircraft determinations of $X_{CO_2}$ gives a correction factor of $0.988 \pm 0.001$ (mean ± standard deviation of the ratios of FTS to aircraft $X_{CO_2}$). Aircraft campaigns over the Park Falls observatory resulted in correction factors of $0.991 \pm 0.002$ calculated in the same fashion as described above.

Using all aircraft overflights from both sites, the average correction factor is $0.990 \pm 0.002$. The FTS $X_{CO_2}$ is therefore biased low by 1.0%. In the $X_{CO_2}$ retrievals, the retrieved $O_2$ column is high by 2.3% (averaged over 260,000 spectra from Darwin and Park Falls), meaning the $CO_2$ column retrievals are approximately 1.3% higher than the in situ scale. This indicates an improvement with the updated spectroscopic line parameters relative to those used for the previous Park Falls analysis (Washenfelder et al., 2006). The improvement highlights the importance of line parameters to the accuracy of the retrieved FTS columns. Figure 7 shows all Park Falls and Darwin FTS-aircraft comparison points, and the best fit line illustrating the average TCCON correction factor. All points overlap with the best fit line within the calculated uncertainties, and the correction factors for Park Falls and Darwin agree with the average correction factor within uncertainties.

The method of calibration described here, and by Washenfelder et al. (2006) places the FTS measurements onto the same scale as the global in situ network. The calibration against aircraft profiles provides a transfer standard between TCCON sites, which will be valuable for future sites joining the network, especially those over which aircraft profiles will not be attained. Ideally timeseries of profiles would be obtained over all sites in order to correct for site-specific and time-varying instrument effects.

8 Error discussion

In integrating the aircraft-based in situ measurements to create a calibrated column measurement we have made assumptions about approximately 20% of the profile. With the 2 µmol mol$^{-1}$ errors attributed to these points, they contribute more than 50% (0.4 out of 0.7 µmol mol$^{-1}$) of the error assigned to the integrated aircraft column. So the missing sections of the in situ profile (especially above the 139 hPa aircraft ceiling) dominate the derived column error. Table 2 summarizes the uncertainty in the column integrated $X_{CO_2}$ over a variety of altitude sampling ranges compared to that obtained here. We now consider the effect of making different assumptions in assigning the upper and lower regions of the profile.

For the upper troposphere and stratosphere we compare the model results with two alternatives. Firstly, we take the alternative model results of Waugh et al. (1997) for the stratosphere, which showed a good match to previous aircraft measurements. The relationship between $CO_2$ and potential temperature in this modeled profile for the tropics during March–April is estimated to be $-0.0125 \mu mol mol^{-1} K^{-1}$, and this gradient applied from the tropopause height upwards. The difference between the two integrated stratospheric profiles is 0.05 µmol mol$^{-1}$, considerably less than the assumed error of 0.18 µmol mol$^{-1}$ for this section of the profile.

We also compare our assumed upper profile with a composite case, consisting of aircraft profile measurements simultaneously acquired over Costa Rica during Costa Rica – Aura Validation Experiment (CR-AVE), which have shown a good match to the profiles during the TWP-ICE campaign from the TTL upwards (Park et al., 2007). The average profile measurement from this campaign is appended to the top of the profiles (above 360 K) used here, and then the original model assumption used to extend the stratospheric profile above the aircraft ceiling. The difference to the total integrated aircraft column dry-air mole fraction

| Table 1. Summary of the column average dry-air mole fractions obtained during the intercomparison between the FTS and the in situ instrument on board Proteus. |
| --- |
| Profile | Aircraft $X_{CO_2}$ | FTS $X_{CO_2}$ | Number FTS spectra |
| --- | --- | --- | --- |
| Upward | 378.1 ± 0.7 | 373.7 ± 0.4 | 49 |
| Downward | 378.3 ± 0.7 | 373.6 ± 0.3 | 37 |

| Table 2. Uncertainties in aircraft integrated columns sampling over specific altitude ranges. |
| --- | --- |
| Sampling range | Uncertainty in integrated $X_{CO_2}$ (µmol mol$^{-1}$) |
| --- | --- |
| 0–4 km | 1.3 |
| 0–12 km | 0.5 |
| These flights | 0.7 |
| 0.3–14.7 km (these flights w/no missing data) | 0.4 |
for both profiles is 0.009 µmol mol\(^{-1}\), and assuming an error of ±0.1 µmol mol\(^{-1}\) for the WB57 values decreases the error estimate by ∼25%. The good agreement with the assumed model profile is heartening, and suggests that the ±2 µmol mol\(^{-1}\) uncertainties attributed are indeed generous, and the assumptions are reasonable.

For the FTS \(X_{CO_2}\) columns, the magnitude of random errors is shown by the clear sky reproducibility of 0.1%. Due to spectroscopic errors, a systematic offset of 1% is introduced to the \(X_{CO_2}\) calculated from FTS measurements of vertical column \(CO_2\) and \(O_2\). These systematic errors are removed via correction to place the FTS determined \(X_{CO_2}\) on the WMO scale for low airmasses. Current spectroscopic linelists contain errors that cause apparent solar zenith angle dependence in the retrieved \(X_{CO_2}\) values; hence we apply an airmass-dependent correction to \(X_{CO_2}\) derived from data at Darwin and Park Falls. These errors may be due to missing lines, or in the line widths. We are confident that with several new spectroscopic measurements becoming available (Robichaud et al., 2008), the systematic errors in the \(O_2\) and \(CO_2\) columns and airmass dependencies can be significantly reduced.

9 Summary and conclusions

An automated solar observatory was deployed to Darwin, Australia, and has acquired near-infrared solar spectra since August 2005. The column \(CO_2\) retrievals were calibrated using aircraft profiles collected over the site during the TWP-ICE campaign in January 2006. These profiles show the \(X_{CO_2}\) determined from FTS vertical column retrievals is underestimated by a factor of 0.988 ± 0.001, based on spectral fits using a modified HITRAN linelist. This bias is in reasonable agreement with that measured at Park Falls (0.991 ± 0.002), suggesting a common cause. The biases in the \(O_2\) and \(CO_2\) column retrievals result from inaccuracies in the spectroscopic line parameter database and the spectral fitting process.

The calibration described here presents a means of placing FTS column-average mole fractions from Darwin on to the same scale as in situ measurements. Data acquired from the site will assist in quantification of tropical \(CO_2\) sources and sinks, where measurements are currently scarce. In addition, this site will provide validation of satellite measurements in this crucial region of the globe.

Appendix A

Airmass correction

There appears to be an airmass dependent artifact (ADA) in the retrieved \(X_{CO_2}\) values, not only for Darwin, but for all TCCON sites. The ADA causes \(X_{CO_2}\) retrievals to be approximately 1% larger at noon than sunrise or sunset, even at clean air sites and at times of the year when there is no physical reason for a diurnal change of \(X_{CO_2}\). The ADA is aliased into the diurnal and seasonal cycles, causing systematic differences between sites. The aliasing problem is particularly important for Darwin and other Southern Hemisphere sites, where the seasonal cycle is smaller than in the Northern Hemisphere. The airmass dependency may be attributable to spectroscopic inadequacies, such as missing lines or errors in temperature-dependencies or line widths, or to instrumental problems such as zero level offsets, continuum curvature and ILS errors due to misalignment.

To derive and apply a correction for the ADA, we assume that on any given day, \(X_{CO_2}\) variation symmetrical about noon is an artifact, and anti-symmetric variation is real. This assumption should be reasonable for unpolluted sites, since the natural effect of photosynthesis and respiration is to cause a maximum \(X_{CO_2}\) at sunrise, and minimum at sunset. For each day, we represent the \(X_{CO_2}\) by three terms, a constant noon-time value (\(\hat{X}\)), a component varying symmetrically about noon (\(S(\theta)\)) and a component varying anti-symmetrically about noon (\(A(t)\)). The amplitudes of each term are determined by a least-squares minimization of:

\[
\chi^2 = \sum_i \left( \frac{X_{CO_2,i} - \hat{X}_{CO_2} - a.A(t) - \beta.S(\theta)}{\epsilon_i} \right)^2
\]

where:
\(\hat{X}_{CO_2}\) = the noon \(X_{CO_2}\)
\(X_{CO_2,i} \pm \epsilon_i\) = the retrieved \(X_{CO_2}\)
and uncertainty for the ith spectrum
\( \theta_i \) = the solar zenith angle for spectrum \( i \) (degrees)
\( t_i \) = the time at which the \( i \)th spectrum was acquired (in days)
\( t_{noon} \) = the time at local solar noon (in days)
\( A(t_i) = \sin (2\pi (t_i - t_{noon})) \)
\( S(\theta_i) = \left( \frac{\theta_i + 13}{90 + 13} \right)^3 - \left( \frac{45 + 13}{90 + 13} \right)^3 \)

\( S(\theta_i) \) is defined to be zero at 45\(^\circ\) because this minimizes the absolute size of the correction and makes it so that \( S(\theta_i) \) is more orthogonal to \( \hat{\chi} \). Making the problem better posed mathematically. These basis function forms were determined from \( X_{CO_2} \) results from Darwin, which are the most plentiful of the TCCON sites, and cover the widest range of solar zenith angles. The decision was guided by the \( \chi^2 \) of minimization. The \( X_{CO_2} \) can then be corrected using Eq. (A2).

\[
X_{CO_2}^{corrected} = X_{CO_2} - \beta \cdot S(\theta_i)
\]  

An average value of beta of \(-0.0075\) is determined using data from Park Falls and Darwin. Figure A1 shows the corrected (grey) and uncorrected (black) \( X_{CO_2} \) and \( X_{CO_2} \) anomaly as a function of solar zenith angle for Darwin and Park Falls. This value of beta results in a correction of \(-0.13\% \) to \( X_{CO_2} \) at 0\(^\circ\) solar zenith angle and \(0.42\% \) to \( X_{CO_2} \) at 80\(^\circ\) solar zenith angle. The between-site and time-based variability in beta is around 10\% of this figure, meaning that the potential bias introduced by applying the airmass correction is less than 0.05\% in \( X_{CO_2} \), significantly below the measurement reproducibility. The correction removes a large portion of the ADA. The consistency of the airmass-dependent correction between sites suggests that the spectroscopic factors contributing to the airmass-dependence are larger than the instrument specific factors.

Acknowledgements. We thank Rex Pearson, John Glowacki, Troy Culgan, Maciej Ryczek and Krzysztof Krzton for the maintenance of the solar FTS. We also gratefully acknowledge comments on the manuscript made by Janina Messerschmidt. Nicholas Deutscher is supported by an Australian Postgraduate Industry Award. The research described in this paper was performed for the Orbiting Carbon Observatory Project at the Jet Propulsion Laboratory, California Institute of Technology, under a contract with the National Aeronautics and Space Administration. TCCON is funded by the NASA terrestrial carbon cycle program, grant NNX08AI86G. This research is also assisted by Australian Research Council Projects DP0879468 and LP0562346 with the Australian Greenhouse Office.

Edited by: I. Aben

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