Structures in the microwave background radiation

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We compare the actual WMAP maps with artificial, purely statistical maps of the same harmonic content to argue that there are, with confidence level 99.7\%, ring-type structures in the observed cosmic microwave background.

1. Introduction

Since its discovery the cosmic microwave background radiation (CMB) has been the subject of very intensive research. The CMB anisotropy provides invaluable information not only about the Universe in the epoch of last scattering but also about the earlier epochs. It follows from the lucky coincidence that before the last scattering the Universe was radiation dominated, and in the presence of radiation pressure it could not evolve and all the structures essentially remained frozen. Therefore although the temperature at the last scattering was rather low (\textasciitilde 1 eV), searching for structures of different types in the CMB anisotropy can possibly reveal facts about radiation (CMB) has been the subject of very intensive research.

2. Creation of artificial maps

Our algorithm for obtaining the maps with artificial CMB temperature distribution was as follows:

1. we choose a maximal multipole number \( L \) (in the present paper we use \( L = 1000 \))

2. for each multipole number \( l \), \( 0 \leq l \leq L \), we choose a suitable number \( N_l \), a positive integer, which is the number of spherical harmonics contributing to the temperature for this given \( l \)

3. for all integers \( k_l, 1 \leq k_l \leq N_l \), we assign a direction \((\theta_{k_l}, \phi_{k_l}) = (\arccos(\pi x_{k_l}), 2\pi y_{k_l})\) where \( x_{k_l} \) and \( y_{k_l} \) are random numbers from the uniform distribution in the interval [0, 1]

4. we take the coefficient \( C_l \) from the known real WMAP spectrum, and define a function \( T_l = T_l(\theta, \phi) \) on the sphere by:

\[
T_l(\theta, \phi) = \sum_{k_l=1}^{N_l} \sqrt{\frac{2l}{N_l}} P_l(\cos(\omega(\theta, \phi, \theta_{k_l}, \phi_{k_l})))
\]

where \( \omega(\theta, \phi, \theta_{k_l}, \phi_{k_l}) \) is the spherical angle between a given point \((\theta, \phi)\) on the sphere and the direction \((\theta_{k_l}, \phi_{k_l})\)

5. we repeat points 2)-4) for each \( l \), \( 0 \leq l \leq L \)

The final formula for the temperature on the artificial map is:

\[
T(\theta, \phi) = \sum_{l=0}^{L} T_l(\theta, \phi).
\]

It depends on our choice of the integers \( L \) and the \( N_l \)’s. There is no preferred direction on the sky since for each \( l \) we choose randomly \( N_l \) directions \((\theta_{k_l}, \phi_{k_l})\) (therefore this procedure has the advantage over averaging over \( m \) in the standard approaches). If we now perform the usual harmonic analysis of this map we get a set of \( C_l \)’s that differ from the WMAP \( C_l \)’s used in the procedure but the difference can be made as small as we wish with the choice of sufficiently large \( N_l \)’s. The artificial maps produced in this way are purely statistical, Gaussian and, depending on the chosen \( N_l \), can reproduce any prescribed spectrum of \( C_l \) with arbitrary accuracy.

Since the spectrum of \( C_l \) obtained from the real maps is (especially for small \( l \)) not precisely known we have produced artificial maps along the average, the upper and the lower \( C_l \) curve. For all of them we obtained essentially the same result.

3. Looking for ring-type structures

We have excluded from all the maps the Milky Way belt (\( \pm 0.4 \) radians above and below the Galactic equator) to avoid introducing structures caused either by the
for any given real or artificial map consisted in the following:

1. a grid of points with HEALPix (k=6) parametrization spreading over the entire sphere has been created;
2. a function \( \Theta_{\epsilon, \gamma_k}(\gamma) \):
   \[
   \Theta_{\epsilon, \gamma_k} = \begin{cases} 
   -\epsilon(\cos(\gamma_k - \epsilon) - \cos \gamma_k)^{-1} & \gamma_k - \epsilon < \gamma < \gamma_k \\
   \epsilon(\cos(\gamma_k) - \cos(\gamma_k + \epsilon))^{-1} & \gamma_k < \gamma < \gamma_k + \epsilon \\
   0 & \text{otherwise}
   \end{cases}
   \]
was defined (this form is dictated by the area measure on the sphere and the condition that a constant temperature should give no contribution to the integral below);
3. for \( N \) (\( N \approx 7000 \)) directions \((\theta^i, \phi^i)\) \((i = 1, \ldots, N)\) belonging to the grid, for all widths of rings \( \epsilon \) \((\epsilon = 0.02, 0.04 \text{ or } 0.08 \text{ radians})\) and for all angular “radii” \( \gamma_k = 0.07 + k \cdot 0.01 \ (k = 1, \ldots, 23) \) an integral
   \[
   I_{(\epsilon, \gamma_k)}(\theta^i, \phi^i) = \int d\Omega' \Theta_{\epsilon, \gamma_k}(\omega(\theta^i, \phi^i, \theta', \phi')) T(\theta^i, \phi^i)
   \]
was calculated \((\omega(\theta^i, \phi^i, \theta', \phi')\) is the spherical angle between \((\theta^i, \phi^i)\) and \((\theta', \phi')\). In this way for a given \( \epsilon \) we get a set of \( N \cdot 23 \approx 160000 \) values \( I_{(\epsilon, \gamma_k)}(\theta^i, \phi^i) \).
4. for a given \( \epsilon \) we order these \( N \cdot 23 \) values of \( I_{(\epsilon, \gamma_k)}(\theta^i, \phi^i) \) (separately for positive and negative values), group them in a number \( n \) \((n = 1500)\) of bins with centers \( x_i \) and for all artificial maps we form the normalized cumulative distribution \( F_n(x_i) \) \((n \text{ stands for the map number and } i \text{ stands for the bin number})\);
5. at each \( x_i \) we average over all \( F_m(x_i) \) to arrive at the average distribution \( F(x_i) \);
6. we use the procedure for comparison of cumulative distribution functions introduced in [4]: for all artificial and 3 real maps W1 – W3 we calculate
   \[
   A = -\frac{a}{N} \sum_{i=1}^{n} d_i \ln(1 - F^n(x_i)) \tag{2}
   \]
where \( a \) is a positive number (we used several numbers, all very large, much larger than \( n \)) and \( d_i \) is the number of points in the \( i \)-th bin. Then, separately for positive and negative distributions, we check the hypothesis that a given map belongs to the distribution \( F(x) \) using the formula given in [4] – the formula quoted there is suitable since it allows for quantitative comparison of distribution functions differing mostly at the tails (as is the case here).

The shape of the function \( \Theta_{\epsilon, \gamma_k}(\theta^i, \phi^i) \) defines what we mean by the ring-type structure – the maximal integral is when the temperature is of the same shape. Since we gather in the same set all values of \( \gamma_k \) the same direction can give several large values of the integral so these values can be correlated. The division into smaller number of bins \( n \) \((n \ll N \cdot 23)\) is intended to partially compensate for these correlations.

4. Results

It is commonly taken for granted (with the notable exception of [3]) that the Cosmic Microwave Background is purely statistical being produced by the quantum fluctuations usually assumed to have taken place during inflation (as the solution in De Sitter space suggests). Therefore it was very unexpected for us to find significant differences (with confidence level 99.7\%\) between the WMAP results and artificial maps (with the same statistical properties as WMAP) that we have created. We believe that the procedure of creating artificial maps described above, being different from the one usually used, is interesting on its own. It gives purely statistical gaussian maps, and reproduces the prescribed harmonic content with the desired accuracy.

We have produced many artificial maps (to generate an artificial map with our procedure is fast) but the procedure for looking for ring-type structure required the number of integrals for any given map to be very large \((N \cdot 23 \approx 160000)\) and the CPU time needed was the main reason to limit ourselves in the present paper to 100 artificial maps only.

The differences between real and artificial maps were both qualitative and quantitative. First of all, for \( \epsilon = 0.08 \) on frequency subband maps W1 – W3 we find two directions around which there is a significant concentration of circular structures distinguished by the large values of \( I_{(\epsilon, \gamma_k)}(\theta, \phi) \). The Galactic coordinates of these two directions are approximately \( (\tilde{\theta}_1, \tilde{\phi}_1) = (2.6, 3.7) \) (which correlates with the so cold “cold spot” [2, 7]) and \((\tilde{\theta}_2, \tilde{\phi}_2) = (2.6, 2.9) \) (this is new and has opposite values of the integral – the highest values from this direction are shown on Fig. 2). We interpret these two directions as two centers of circular rings in the temperature distribution. The reason for this interpretation is as follows: in a close vicinity (about \( 8 \cdot 10^{-3} \text{ sr} \)) of the first of these directions, there are about 80 centers of circular structures (and there are none outside). A similar situation occurs around the second direction – there we find about 40 centers in about the same solid angle as in the first case. We also found many more potential centers of ring-type structures but since they were less significant statistically we do not include them in the analysis. It is perhaps worthwhile to note that the positions of both of the main two ring-type structures and the less significant ones correlate approximately with the positions of the structures found in [3].

The quantitative comparison between the real W1 – W3 maps with the artificial maps regarding the existence
of structures with $\epsilon = 0.08$ yields the following:

For the procedure described above we get for the real maps $W1 - W3$ with $n = 1500$, $a/n = 2000$ the values $A = 107.9$, $133.5$, $84.8$ respectively, while only 6 out of 100 artificial maps have $A > 1$ (the largest value among these 6 maps is $A \approx 52$, the other 5 have much smaller $A$). According to [4] the probability of $A$ bigger than $\sigma$ for large $a/n$ is approximately equal to $(1 \ll \sigma \ll a/n)$

$$P(A > \sigma) = \frac{n}{a} \ln \frac{a}{n\sigma}$$

(3)

so we get for the real map $W2$ the probability of a statistical fluctuation $\approx 0.3\%$, which means that with 99.7% confidence level we can reject the hypothesis that the values of the integrals on the real maps are purely statistical and state that the ring type structures on the background radiation maps are not statistical fluctuations. In favor of this statement we can also observe that for the value of $I(\epsilon,\gamma_0)(\theta, \phi)$ greater than 3 on the map of frequency band $W2$ only 81 points are in the northern hemisphere and 1430 in the southern one – the difference is similar for two other spectral types but there are no such differences for the artificial maps.

It is interesting to note that for $\epsilon = 0.02$ and $\epsilon = 0.04$ the differences between the real and artificial maps are much smaller.

For comparison we have also imposed a condition that around the same direction ($\theta, \phi$) we have both large positive and large negative integrals. It turned out that none of the real maps shows any statistically significant presence of such structures – it points to the presence of one sharp circular edge in the temperature distribution and not two or more (which may be there but milder).

In figure 1 we have plotted the directions of the suspected ring-type structures on the real WMAP maps in 3 frequency bands $W1 - W3$ (with the galactic disk excluded, as described in the text). In fig 2. we have plotted the upper part of the histogram of the cumulative distribution functions for all artificial maps and three real maps $W1 - W3$ (for positive value of the integral).

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Fig.1. Position of centers of circular structures on the real maps $W1 - W3$
Fig. 2. The tail of the distributions of the integrals of the 100 artificial (blue) and 3 real (red) maps