On explosive boiling of a multicomponent Leidenfrost drop

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The gasification of multicomponent fuel drops is relevant in various energy-related technologies. An interesting phenomenon associated with this process is the self-induced explosion of the drop, producing a multitude of smaller secondary droplets, which promotes overall fuel atomization and, consequently, improves the combustion efficiency and reduces emissions of liquid-fueled engines. Here, we study a unique explosive gasification process of a tricomponent droplet consisting of water, ethanol, and oil (“ouzo”), by high-speed monitoring of the entire gasification event taking place in the well-controlled, levitated Leidenfrost state over a superheated plate. It is observed that the preferential evaporation of the most volatile component, ethanol, triggers nucleation of the oil microdroplets/nanodroplets in the remaining drop, which, consequently, becomes an opaque oil-in-water microemulsion. The tiny oil droplets subsequently coalesce into a large one, which, in turn, wraps around the remnant water. Because of the encapsulating oil layer, the droplet can no longer produce enough vapor for its levitation, and, thus, falls and contacts the superheated surface. The direct thermal contact leads to vapor bubble formation inside the drop and consequently drop explosion in the final stage.

多组分微滴在热表面的爆裂沸煮

多组分微滴爆裂沸煮的过程对于改善发动机的燃烧效率和减少排放具有重要意义。通过高速摄像监测整个气化过程，我们发现了一种独特的爆炸气化现象。在这个过程中，挥发性最强烈的组分——乙醇——优先蒸发，触发了油微滴/纳米滴在剩余滴中的核化，最终形成了一个不透明的油包水微乳液。这些微滴随后合并成一个大滴，包裹了剩余滴中的水。由于被油层包覆，这个滴不能再产生足够的蒸汽来使其浮起，从而落下来并接触热表面。直接的热接触导致了在滴内部的气泡形成，进而引发整个滴的爆炸。

**Significance**

加热、蒸发和混合滴通过喷射是一个内在的热能生产过程。通过自我诱导的微滴气化，可以实现对多组分燃料混合物的气化，而这种混合物具有优化的挥发性和相容性，为改进气化特性提供了潜在的途径。这项工作全面地探究了气化过程中的三相滴，它最初是在Leidenfrost状态下的一个超热表面，从初始的三相气化开始，最终通过自我封装和微爆裂来实现整个滴的爆炸。这项工作充分展示了现象的丰富性，这种现象可能在多组分滴系统中出现，对相转变的影响。

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surface at 400 °C. As expected, the ouzo drop was in the levitated Leidenfrost state (Fig. 1A), and a vapor layer was formed below it, isolating it from direct contact with the solid surface and, thus, constituting the so-called film boiling state. A slightly curved quartz lens was used as the solid surface to prevent the drop from moving around, as detailed in Materials and Methods.

The observed boiling process takes more than 20 s, with the evaporating ouzo droplet experiencing four distinct stages. Fig. 1 shows representative experimental snapshots of the drop in each stage from a side-view recording, together with the corresponding sketches in the cross-section view. As shown in the first image of Fig. 1A, the ouzo drop is initially miscible and, hence, optically transparent. Subsequently, the preferential evaporation of the most volatile ethanol component reduces the oil solubility. Consequently, spontaneous nucleation of the oil microdroplets/nanodroplets (arrows in the second image in Fig. 1A)—i.e., the ouzo effect—is triggered, as was similarly observed by Tan et al. (22, 29) for an evaporating sessile ouzo droplet, with the drop sitting on a solid substrate with a three-phase contact line.

This marks the end of the first stage. Many microdroplets appear and are advected around by the internal flow (the third image in Fig. 1A). Then, more and more oil microdroplets visibly appear in the drop—at the beginning, only on the drop surface, but later, also in the bulk of the drop. After about 1.28 s, the drop is full of the oil microdroplets and becomes opaque (the fourth image in Fig. 1A), defined as the end of stage 2. This is, again, similarly seen by Tan et al. (22, 29) for an evaporating ouzo droplet on a solid substrate. The entire process of the first two stages is sketched in Fig. 1D. Starting from stage 3, the oil microdroplets start to coalesce and form large and transparent oil droplets, as shown in the first image of Fig. 1B. These gradually merge into one entity, which constitutes part of the drop surface. Finally, an oil cap is formed on the upper half of the clear aqueous drop, as shown in the second image of Fig. 1B. Remarkably, the oil cap then gradually spreads over the drop surface (the third image in Fig. 1B), as also seen for a drop on an atmospheric surface (30, 31). The process of the third stage is sketched in Fig. 1E. After the wrapping-up process is finished, a fully oil-capsulated Leidenfrost drop has been established. However, this aqueous-in-oil drop cannot be stably levitated, due to the lack of sufficient vapor, and leads to the last stage of the drop’s life, as shown in Fig. 1C. Specifically, in the first image in Fig. 1C, the top part of the drop appears blurry, due to the unstable motion of the drop surface. With the levitating vapor layer formed by the vaporizing volatile water being cut off by the encapsulating surface oil layer, the droplet falls down and contacts the superheated surface. The droplet is now heated rapidly, eventually leading to the nucleation of the encapsulated volatile liquid core and, hence, the explosion of the drop as the end of its life, as shown with the third and fourth images of Fig. 1C and the sketches in Fig. 1F.

To start a quantitative analysis, we calculate the volume \( V \) for the evolution of the Leidenfrost ouzo drop, assuming that the drop maintains its axisymmetric shape during the entire evaporation process. As shown in Fig. 2A, the drop size decreases rapidly within 24 s. Blue circles show experimental results. The scatter of the data originates from the drop deformation. The red dashed line and the blue area, respectively, show the fitting line and the changing range of the volume. It is seen that during stage 4, the drop is unsteady and asymmetric. Thus, it is hard to determine the exact volume evolution in time, and only the drop volume before the explosion moment is calculated. The red circle with the error bar shows the average volume in

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**Fig. 1.** The entire boiling process of an ouzo drop initially in the Leidenfrost state can be divided into four stages. For each row, A–C are experimental results, and D–F are the corresponding sketches in the cross-section view. In the sketches, the orange parts represent the oil-rich phase and the blue parts the water-rich phase. \( t_1 \), \( t_2 \), \( t_3 \), and \( t_4 \) indicate the end moment of each stage. (A and D) Stages 1 and 2: The drop transforms from the transparent state to the early milky state. Oil microdroplets/nanodroplets first nucleate on the interface (see arrows), marking the end moment \( t_1 \) of stage 1, defined as \( t_1 = 0 \). Then, the drop is filled with oil emulsions and, accordingly, becomes fully opaque at \( t_2 \). (B and E) Stage 3: Tiny oil microdroplets gradually coalesce and, finally, form an oil cap on the upper half of the drop. The drop transforms from the early milky state to a transparent state with two separate phases. The oil cap tries to cover the drop and, finally, wraps the drop at \( t_3 \). (C and F) Stage 4: The drop, which previously was in a stable film-boiling Leidenfrost state, now becomes unstable. The vapor generated from the oil encapsulation is no longer sufficient to levitate the drop, and, consequently, the drop directly contacts the solid surface and explodes at \( t_4 \). (Scale bars: 2 mm.)
stage 4. To gain an overall view of the evaporation rate, we calculate the absolute value of the volume decreasing rate per unit area, |\( \dot{V} \)| = |\( \frac{dV}{dt} \)| / \( S_o \), where \( S_o \) is the entire surface area of the drop determined from the side profile of the drop. As shown in Fig. 2B, the red dashed line is calculated by the fitting lines of volume \( V \) and entire area \( S_o \) of the drop. The blue area shows the possible range of |\( \dot{V} \)|. The red circle with the error bar shows the average |\( \dot{V} \)| in stage 4. Moments of \( t_1 \), \( t_2 \), \( t_3 \), and \( t_4 \) are marked through the vertical dashed lines.

Fig. 2. Quantitative analysis of the entire evaporation process. (A) Volume \( V \) versus time \( t \) of the ouzo drop. Blue circles show experimental results. The spreading of the data originates from the drop deformation, leading to an error in identifying the droplet volume from the cross-section (shown in the blue area). The red dashed line is the fitting line of the volume. The red circle with the error bar shows the average volume in stage 4. (B) Evolution of the absolute value of the evaporation volume flux rate per unit area |\( \dot{V} \)| versus time. The red dashed line is calculated by the fitting lines of volume \( V \) and entire area \( S_o \) of the drop. The blue area shows the possible range of |\( \dot{V} \)|. The red circle with the error bar shows the average |\( \dot{V} \)| in stage 4. Moments of \( t_1 \), \( t_2 \), \( t_3 \), and \( t_4 \) are marked through the vertical dashed lines.

decreases rapidly, as shown with the red circle in Fig. 2B. The corresponding ending moments of each stage \( t_1 \), \( t_2 \), \( t_3 \), and \( t_4 \) are marked by the dashed vertical lines in Fig. 2. Three further and independent experimental realizations of the same experiment are shown in SI Appendix, all giving similar results.

As discussed above, the Leidenfrost ouzo drop experiences complex boiling dynamics, including four different stages, which we now discuss in more detail.

Stages 1 and 2: From Transparent Miscible Ouzo Drop to Emulsification due to Microdroplet Nucleation. The first characteristic event during the boiling process of the Leidenfrost ouzo drop is the transformation from the initial miscible drop to an opaque one due to the nucleated oil microdroplets, i.e., an emulsified drop. The emulsification is due to the higher evaporation rates of ethanol, which leads to lower oil solubility and, eventually, nucleation of oil microdroplets (22).

The temperature of the Leidenfrost drop can be approximated by the boiling temperature of the mixture (32). At the initial drop composition, the boiling temperature is around 80 °C (33), which is higher than the boiling point of pure ethanol (78 °C), but lower than that of pure water (100 °C), and much lower than that of transanethole (234 °C). Consequently, ethanol preferentially vaporizes, and its reduced concentration leads to the corresponding reduction of the oil solubility, which, in turn, induces the ouzo effect (22), characterized by the spontaneous nucleation of oil emulsion. Fig. 3A gives four snapshots, showing that the oil microdroplets start nucleating near the drop surface and then are advected by the convection inside the drop.

We used the microdroplets as tracer particles to reveal the flow field in the drop by vertically shining a laser sheet at the center of the drop. A particle-image velocimetry (PIV) calculation (PIVlabotary) shown in Fig. 3B displays the motion of the microdroplets, which move upward along the surface and then downward inside the drop (Movie S1). Two large vortices form in the cross-section, corresponding to one toroidal vortex in the three-dimensional (3D) drop. For a gravity-flattened drop (see the third image in Fig. 3A), the liquid at the drop base is entrained by the viscous flow in the vapor layer from the base center to the periphery, which contributes to the flow moving upward near the interface and downward near the axis (34).

In addition, the temperature is lower at the apex of a large Leidenfrost drop (34), again giving a higher surface tension there. The temperature gradient also causes thermal Marangoni flows along with the interface from the base to the apex (34). However, a higher temperature at the drop base may induce a higher evaporation rate of ethanol, giving a higher surface tension there. The concentration gradient may induce solutal Marangoni flows along with the interface in the opposite direction of the thermal Marangoni flow. Apparently, the measured motion of the microdroplets (Fig. 3B) suggests that solutal Marangoni flows are weaker in the current situation.

Stage 3: Water-in-Oil Drop by Oil Wrapping-Up. The opaque emulsified drop is in a transient state, as the nucleated oil microdroplets are metastable. The continuing oil nucleation and the microdroplet coalescence lead to the growth of the microdroplets. However, the larger oil microdroplets cannot follow the flow closely. As shown in Fig. 4A, because of buoyancy, they tend to float around the apex of the drop, and, consequently, the oil phase merges at the top of the drop.

To provide some insight into the merging and floating processes of the oil microdroplets, we considered the motion of the microdroplets that is controlled by two forces—i.e., the drag force \( F_d \propto \mu R_o v_o \) and the buoyancy force \( F_b \propto \Delta \rho R_o^3 v_o^2 \), where \( \mu \) is the liquid dynamic viscosity in the drop, and \( R_o \) and \( v_o \) are the radius and velocity of the microdroplets. Thus, the drag-buoyancy force ratio is inversely proportional to the droplet size.

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squared, \( \frac{F_d}{F_b} \propto \mu \nu_0 \Delta \rho^{-1} R_{eq}^{-2} \). When the microdroplets are small, the drag force prevails, and, hence, they follow the flow closely. However, the buoyancy force starts to dominate for the larger microdroplets. Fig. 3C shows the temporal evolution of the velocity at the half-height of the drop. For all three times, the maximum velocity is around the center of the drop. At 8.45 s, the maximum magnitude is about 15 cm/s. Recognizing that, since most ethanol may be depleted at this time, to evaluate the dimensionless parameter of the water-rich aqueous drop, we can use the physical properties of liquid water at the boiling temperature. The Péclet number, representing the ratio of convective transfer over the diffusive one, is \( Pe = \frac{v_{\text{max}} R}{\alpha} \approx 1741 \), where \( v_{\text{max}} \) is the maximum velocity within the drop, \( R \) is the equivalent radius of the drop, and \( \alpha \) is the thermal diffusivity of liquid water. Here, we only consider the thermal transport, as the thermal effect is the dominant one in the current stage. The advective transport rate is much higher than that of thermal diffusion, suggesting that the diffusion is way too slow to level out temperature.

\[ \text{Fig. 3.} \quad \text{From transparent miscible ouzo drop to emulsification due to microdroplet nucleation.} \quad (A) \quad \text{The white nucleated oil areas first occur on the interface and are advected by the internal flow. A laser sheet vertically shines through the center of the drop from the top for visualization. In the cross-section of the drop, the drop is transparent in the beginning. Then, the nucleated oil microdrops are distributed through two large vortexes. Subsequently, the oil microdroplets distribute throughout the entire drop. As these oil microdroplets coalesce, the drop becomes more transparent, and a strong flow is observed within the drop.} \quad (B \text{ and } C) \quad \text{The images of } A \text{ are used to analyze the velocity within the drop using PIV.} \quad (B) \quad \text{The internal velocity distribution in the cross-section of the drop shows that the flow moves downward near the axis and upward near the interface. Two large vortices form in the cross-section, corresponding to one toroidal vortex in the 3D drop. For all three times, the maximum velocity is around the center of the drop. At 8.45 s, the maximum magnitude is about 15 cm/s.} \quad (C) \quad \text{Temporal evolution of the velocity at the half-height of the drop. With advancing time } t, \text{ both the maximum velocity and the mean velocity decrease.} \]

\[ \text{Fig. 4.} \quad \text{Stage 3 in the drop history.} \quad (A) \quad \text{The internal oil microdroplets coalesce, and, correspondingly, the mean size of oil droplets increases with time. The drop becomes more and more transparent.} \quad (B) \quad \text{The oil cap gradually covers the aqueous drop and, finally, wraps around it.} \quad (C) \quad \text{The corresponding sketches of the wrapping process in the side view. Note that water is trapped inside the droplet. (Scale bars: 1 mm.)} \]
differences. Furthermore, both the maximum velocity and mean velocity decrease with time. Since the small microdroplets grow and the average driven flow slows down, more and more oil droplets float up and accumulate at the top. As demonstrated in Fig. 4A, the oil droplets float around the apex of the drop, and then gradually merge into an oil cap. The phenomenon is new as compared to the evaporation of an ouzo drop directly in contact with a solid surface, as there, the contact line has major effects on the evaporation process (22, 29). In such a case, the oil droplets accumulate at the contact line and form an oil ring. In contrast, here, the ternary Leidenfrost drop does not have a contact line, and the oil droplets can move freely within the Leidenfrost drop. Oil droplets start to coalesce and gradually merge into one entity. Finally, the water-rich aqueous drop is attached to an oil cap.

The reduction of the ethanol concentration leads to an increase of the drop surface energy and, correspondingly, changes the wetting condition of the oil cap on the drop. Recognizing that the surface tension of water is very high, the oil cap tries to cover the aqueous drop to reduce the surface energy. Fig. 4B then shows a series of snapshots that capture the spreading of the oil cap on the drop surface. Additionally, the upward interfacial flow competes with the spreading of the oil from the top, which causes repeated advancing and retreating spreading of the oil film (Movie S2). In the end, the oil film spreads and wraps the entire drop. The sketches in Fig. 4 show the wrapping process of the oil film, as discussed. Consequently, the emulsified drop completes the oil encapsulation and becomes a water-in-oil Leidenfrost drop. We note in passing that a similar wrapping-up phenomenon has also been observed in other fluids systems by Li et al. (30), Tan et al. (31), and Kadota et al. (21).

**Stage 4: The Final Drop Explosion.** According to the lubrication approximation, the lifting force induced by the vapor layer should balance the weight of the drop. The thickness of the vapor layer can then be evaluated as (35)

\[
h = \left( \frac{3\pi \mu_v |\nabla| \rho_w r_b^2}{2 - \rho_c G^2} \right)^{1/3}, \tag{1}
\]

where \( r_b \) is the horizontal extent of the drop bottom surface, and \( \mu_v, \rho_v, \) and \( \rho_c \) are the dynamic viscosity of vapor, the vapor density, and the liquid-phase density of the evaporating component (water in stage 3 and oil in stage 4), respectively. The physical properties of water and oil are given in SI Appendix. Because the volume of oil is much smaller than that of water, and the density difference between oil and water is small, the gravity of the drop can be calculated by \( G = \rho_{l,w} V_g \), where \( \rho_{l,w} \) is the density of liquid water, and \( g \) is the gravitational acceleration.

As shown in Fig. 5A, the vapor-layer thickness near the end of the last stage 3 is above 30 \( \mu \)m when water is the main evaporation phase. However, the thickness decreases to around 10 \( \mu \)m based on Eq. 1 when the drop is wrapped with the oil phase in stage 4, as shown in Fig. 5A. The range of the blue area and the value of the error bar are explained in SI Appendix. The very thin vapor layer results from the extremely slow evaporation rate of the oil phase, as shown in Fig. 2B. Such a thin vapor layer cannot stably levitate a millimeter-sized drop (32). A small perturbation will induce direct contact between the drop and the superheated solid surface, and this local contact consequently triggers microexplosion of the water droplet fully entrapped by the oil layer.

We combine the bottom and side recordings, as presented in Fig. 5B, to identify the moment around the explosion. In the first column, there is no contact yet, as shown in Lower. Then, in the second column, the dark area marked in the red circle demonstrates that the drop contacts the heated substrate. Once the drop contacts the superheated substrate, the temperature of the contacting liquid increases rapidly above the oil boiling point, \( T_c \), which is much higher than that of water, \( T_{w} \). Thus, the water phase next to the oil contacting area is further superheated, and vapor bubbles are observed in the green circle of Fig. 5B. Once the drop is unable to sustain the inner high pressure of the vapor bubbles (36), the drop undergoes a violent explosion (see the blue circle of Fig. 5B).

**Conclusions and Outlook**

In conclusion, we have experimentally investigated the entire gasification process of a multicomponent Leidenfrost drop of ouzo composition. It is found that, upon gasification, the preferential evaporation and, hence, concentration reduction of the most volatile ethanol leads to the massive formation of oil microdroplets, which, in turn, agglomerate and form a surface layer that encapsulates the remaining aqueous drop. The

**Fig. 5.** Stage 4 in the drop history. (A) Calculated vapor-layer thickness (based on Eq. 1) versus time during the end of stage 3 and in stage 4. The final vapor-layer thickness is above 30 \( \mu \)m in stage 3, but it decreases to around 10 \( \mu \)m in stage 4. The red dashed line and the blue area, respectively, show the average value and the changing range of the vapor-layer thickness. The red point with the error bar indicates the average vapor-layer thickness in stage 4. (B) Explosion of the drop. Combining the side (B, Upper) and the bottom (B, Lower) views, after the drop violently oscillates for a while, it locally contacts the heated substrate (shown in the red circle). Vapor bubbles are generated within the drop at the contact area (shown in the green circle). The vapor bubbles formed have large internal pressure, which causes a violent explosion in the end (shown in the blue circle). (Scale bars: 1 mm.)
aqueous drop, within the encapsulating oil layer, is then heated to the limit of superheat and, consequently, nucleates near the superheated surface, hence rupturing the entire droplet.

The found phenomena illustrate the beauty and richness of multicomponent drop systems with phase transition, which are also relevant in many combustion processes.

Materials and Methods

Experimental Setup and Procedures. A high-speed camera (Photron Fastcam NOVA S12) with macro lens (Nikon 105 mm) was placed in the side view, and another high-speed camera (Photron Fastcam AX200) with a long-distance microscope (Navitar) was placed in the bottom view to simultaneously record the entire process. A slightly curved quartz lens was used as a substrate to limit the movement of drops. A sapphire base was placed under the quartz lens to improve the temperature uniformity of the substrate. Both of them were heated by an aluminum block heater. The surface temperature of the substrate, $T_s$, was measured by a thermocouple attached to the surface. The transparent ouzo drop solution was prepared with an initial composition of 31.80% (volume [vol/vol]) ultrapure water, 66.05% (vol/vol) ethanol ($\geq 99.7\%$), and 2.15% (vol/vol) transanethole oil (Sigma-Aldrich; 99%), which located the liquid to be initially in the one-phase regime and can easily enter the nucleation region ($22 \pm 29\%$). A large ouzo drop produced by a microliter-pipette was deposited on the superheated substrate. The initial volume of the drop was 100 $\mu$L. The quantitative calculation started around 50 $\mu$L, and the substrate temperature was around 400 $^\circ$C.

Image Analyses. The images were analyzed by using a Matlab code. The shape of the drop was assumed to be axisymmetric (37), and the calculation on the volume evolution started from stage 2.

Data Availability. All data are included in the article and SI Appendix.

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