Validation of Neutron Evaluated Data Based on The Experimental Reactivity Worth of Tungsten Target in CiADS
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I. INTRODUCTION
Recently, Chinese initiative Accelerator Driven Systems (CiADS) project has been started in China [1]. For CiADS, the primary goal is to achieve the integration of three systems, i.e. a subcritical lead-bismuth reactor, a heavy metal spallation target and an intense-beam proton linac [2]. During the early stages of ADS projects, some experiments [3-7], which coupled the proton accelerator with the spallation target, have been performed by the researchers. However, the experimental studies about the coupling behaviors between the nuclear reactor and the spallation target have been rarely investigated. In order to ultimately build the CiADS experimental setup, it is necessary that the nuclear reactor and spallation target coupling experiments should be performed. The tungsten granular spallation target has been proposed by Institute of Modern Physics, Chinese Academy of Sciences (IMPCSAS) as an innovative spallation target in the world. The tungsten granular target, which is placed in the center of the reactor, is impinged by the high-energy protons to produce spallation neutrons to maintain the stable operation of the system [8]. Meanwhile, the spallation target, as the structural material, has an effect on the reactor physical parameters, such as neutron flux, neutron spectra, system reactivity and so on. The mass fraction of tungsten is 93% in the spallation target, therefore tungsten cross sections in evaluated data libraries can impact directly on the accuracy of reactor physical calculations in CiADS. In the past experiment on tungsten data evaluation, tungsten materials are generally used as reflecting layer or fuel diluent materials and almost placed in fast-spectrum benchmark experiment setups [9-15]. However, the relevant studies on the tungsten spallation materials in thermal-spectrum reactor are still scarce. In this paper, the measurement of the cylindrical tungsten target reactivity worth was performed on VENUS-II light water reactor by the period method. And then the experimental tungsten target reactivity worth was compared with the results calculated by MCNP with the different libraries, i.e. ENDF/B-VII.0 [16], ENDF/B-VII.1 [17], JENDL-4.0 [18], JEFF-3.2 [19] and CENDL-3.1 [20]. The difference between the experimental and simulated results is discussed, which provides some reference for the selection of nuclear data libraries on CiADS neutronics simulations.

II. EXPERIMENT
II.1 Experimental facility
The VENUS-II zero-power facility, which consists of a light water reactor and a lead reactor, firstly reached the criticality on December 2016 through the cooperation of China Institute of Atomic Energy (CIAE) and Institute of Modern Physics, Chinese Academy of Sciences [21]. The structures and materials in the VENUS-II light water reactor are relatively simple, which is helpful in validating the neutron nuclear data libraries. Therefore, the present experiments were carried out on the light water reactor. The physical photos of VENUS-II light water reactor are shown in Fig. 1. Light water was used as the moderator and reflector, in which the reactivity worth has been performed on the light water reactor of CIAE in order to verify the neutron evaluated data related to the engineering design of Chinese initiative Accelerator Driven Systems (CiADS). The reactivity worth of the tungsten target was measured and processed as -1.234±0.114mk by a period method. By comparing the results of experiment and simulation, the simulated results from ENDF/B-VII.0, JENDL-4.0 and JEFF-3.2 are higher than the experimental result, however that from CENDL-3.1 is lower. The result from ENDF/B-VII.1 library shows better agreement with the experiment one and the relative deviation is less than 2%. Through the analysis of the differences of the results, non-tungsten elements cross sections in the ENDF/B-VII.1 mainly affect the tungsten radiation capture and elastic scattering reaction rates in the energy range of 10^{-9}-10^{-7} MeV, which results in a better simulated tungsten target reactivity worth value. Therefore, it is recommended that the tungsten target reactivity worth should be calculated with the ENDF/B-VII.1.

Keywords—Tungsten target reactivity worth, nuclear data library, MCNP, VENUS-II, period method

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\[ \rho = \frac{\Lambda}{T/\ln 2} + \sum_{\alpha=1}^{m} \frac{\beta_{\alpha eff}}{1 + (\lambda T / \ln 2)} \]

Where, \( \rho \) is the reactivity; \( \Lambda \) is the neutron generation time; \( T \) is the nuclear reactor doubling time; \( m \) is the number of effective delayed neutron groups; \( \beta_{\alpha eff} \) and \( \lambda \) are respectively the \( \alpha \)th group effective delayed neutron fraction and the decay constant of the \( \alpha \)th group delayed neutron precursor.

The double period \( T \) values were measured by stopwatch and reactor power monitor [23] in different supercritical states of the reactor, and the supercritical states were maintained by adjusting the fuel rods number in the reactor. During the experiment, the four supercritical states to be measured are as follows: (1) 961 fuel rods without anything in the spallation zone, (2) 960 fuel rods without anything in the spallation zone, (3) 963 fuel rods with the target clad in the spallation zone, and (4) 968 fuel rods with tungsten target (include cylindrical tungsten target and target clad) in the spallation zone. More details on the experiments were reported in the Ref. [24]. Meanwhile, the reactor kinetic parameters (\( \Lambda \), \( \beta_{\alpha eff} \) and \( \lambda \)) were respectively calculated by MCNP6 with different libraries, i.e., ENDF/B-VII.0, ENDF/B-VII.1, JENDL-4.0, CENDL-3.1 and JEFF-3.2. And then the reactivity \( \rho \) values were calculated by substituting the measured \( T \) and the calculated kinetic parameters into the Eq. (1). At last, according to the formula \( \Delta \rho = \rho_{\text{mea}} - \left[ \rho_{\text{N06}} + \frac{5}{\Lambda} (\rho_{\text{N06}} - \rho_{\text{N06}}) \right] \), the tungsten target reactivity worth \( \Delta \rho \) was calculated as \( -1.234 \times 10^{-5} \)mk, where \( \rho_{\text{N06}} \), \( \rho_{\text{N06}} \), \( \rho_{\text{N06}} \) and \( \rho_{\text{N06}} \) represent the measured reactivity values under the above conditions (1)-(4), respectively. The experimental error comes from the deviations of \( T \) measurement and the uncertainties of the simulated kinetic parameters [25].

III. Monte Carlo Simulation

III.1 Simulation method

In this paper, the reactivity worth of the cylindrical tungsten target is simulated by MCNP software, which is produced and developed by Los Alamos National Laboratory. In MCNP simulations, VENUS-II light water reactor model was created accurately at real geometry and material. Meanwhile, the loading number of fuel rods in the light water reactor was set as the number (totally 961 rods) of initial fuel rods in the experiment and kept the constant. And then the simulation process had two steps: (1) The target clad was placed in the spallation target zone; (2) The tungsten target replaced the target clad in the spallation target zone. Therefore, the reactivity worth of the tungsten target was calculated as the difference between the reactor reactivity values with the tungsten target and the target clad in the spallation target. The simulation method can reduce the simulated error of the tungsten target reactivity worth results and satisfy the accuracy demands.

III.2 Details description

The effective multiplication factor \( k_{\text{eff}} \) values of the two steps in Part 3.1 were calculated by KCODE card in MCNP with 3000 cycles and 5 \( \times \) 10^5 histories per cycle, which caused the estimated standard deviation of about 2pcm (1pcm=10^(-5)) for the \( k_{\text{eff}} \) results. The temperature values of the material cross sections were set as the room temperature 293k and remained unchanged. Considering the upward scattering of the thermal neutrons, the thermal neutron scattering library S (\( \alpha, \beta \)) was used to achieve the accurate neutron transportation calculations in the light water reactor. Then, the \( k_{\text{eff}} \) was converted to the reactivity \( \rho \) through the equation \( \rho = (k_{\text{eff}} - 1) / k_{\text{eff}} \). The simulated tungsten target reactivity worth was calculated as the difference between the reactivity values with the tungsten target and the target clad in the spallation zone.

To examine the different neutron nuclear data libraries, the cross sections of all materials in the reactor, i.e. target, fuels, moderator and among others, are respectively from ENDF/B-VII.0, ENDF/B-VII.1, CENDL-3.1, JENDL-4.0 and JEFF-3.2. And then five corresponding tungsten reactivity worth values were obtained and the errors originated from the statistical errors of MCNP simulation.

IV. Results and Discussion

The experimental and calculated results of tungsten target reactivity worth values are shown in Fig. 2, where the shaded area indicates the experimental results within the error range, and the five dots respectively represent the MCNP simulated results with five cross section libraries. In Table I, the calculated results in the second column are respectively simulated by MCNP with the different cross section libraries. Meanwhile, the ratio values of the calculated and experimental results (C/E) are presented in the fourth column. From the obtained results, it can be observed that:

(1). Tungsten target reactivity worth values, obtained by the experiment and simulation, are all negative, which causes the reduction of the reactor reactivity.

(2). Comparing with the experimental result, the simulated results of ENDF/B-VII.0, JENDL-4.0 and JEFF-3.2 are higher, whereas the one of CENDL-3.1 is lower.

(3). The simulated result with ENDF/B-VII.1 has a good agreement with the experiment one and the relative deviation is less than 2%.

![Tungsten target reactivity worth values from the experiment and the simulation.](https://doi.org/10.1051/epjconf/202022504026)

**Table I**

| Libraries      | Tungsten target reactivity worth (mk) | MCNP | Period method | C/E  |
|----------------|--------------------------------------|------|---------------|------|
| ENDF/B-VII.0   | -1.090±0.029                         |      |               | 0.883|
| ENDF/B-VII.1   | -1.255±0.029                         |      |               | 1.017|
| CENDL-3.1      | -1.407±0.029                         |      |               | 1.14 |
| JENDL-4.0      | -1.129±0.029                         |      |               | 0.915|
| JEFF-3.2       | -1.050±0.029                         |      |               | 0.851|
In order to find out the reason of the differences among the tungsten reactivity worth values from five libraries, the affecting factors of tungsten reactivity worth were divided into two aspects: the tungsten element and non-tungsten elements. Meanwhile, the simulation results from ENDF/B-VII.0, JENDL-4.0 and JEFF-3.2 are basically in agreement within the error ranges (see Fig. 2), so that the results are grouped into the same group and one calculated by ENDF/B-VII.0 was chosen to represent the group. Therefore, the simulated result from the ENDF/B-VII.1 is only compared with those from the ENDF/B-VII.0 and CENDL-3.1. First, the reason for the difference between the tungsten reactivity worth results with the ENDF/B-VII.1 and CENDL-3.1 is mainly the discrepancy of the tungsten radiative capture cross sections at the epithermal neutron region [24].

Next, the study focused on analyzing the difference between the tungsten reactivity worth values from ENDF/B-VII.1 and ENDF/B-VII.0. Therefore, the tungsten reactivity worth values in the three cases are calculated and shown in Table II. The table shows the results in the case 1 and case 2 are almost consistent with each other within errors. However, the result in case 1 is remarkably different from that in case 3. It can be concluded that the non-tungsten element cross sections are the main reason for the discrepancy of the results between ENDF/B-VII.0 and ENDF/B-VII.1, whereas the impact of the tungsten element cross sections on the difference is almost negligible.

### Table II

| Case No. | Tungsten cross sections | Non-tungsten cross sections | Target worth (mk) |
|----------|-------------------------|-----------------------------|-------------------|
| 1        | ENDF/B-VII.1            | ENDF/B-VII.0                | -1.110±0.029      |
| 2        | ENDF/B-VII.0            | ENDF/B-VII.0                | -1.090±0.029      |
| 3        | ENDF/B-VII.1            | ENDF/B-VII.1                | -1.255±0.029      |

In the neutron field of the reactor, the reaction channels in the tungsten target mainly include the radiative capture, elastic scattering, inelastic scattering, (N, 2N) and (N, 3N). To analyze the effect of the each reaction channel on the tungsten reactivity worth, the reaction rates of each reaction channel were respectively increased by 1.3 times the original values for the sensitivity analysis. More specifically, the input card in case 3 of Table II was calculated by the perturbation function of MCNP and the variations of the reactivity ρ values are shown in Table III. For the specific reaction type, it can be seen that the radiative capture causes a negative reactivity in the system, whereas the elastic scattering and inelastic scattering respectively cause a positive reactivity. (N, 2N) and (N, 3N) have little impact on the reactivity value. Therefore, it can be concluded that the reaction rates discrepancies of the radiative capture, elastic scattering and inelastic scattering in the tungsten target cause the difference of the tungsten reactivity worth values. Moreover, the reactivity variation caused by the radiative capture is the largest, that is, the radiative capture reaction rates have the greatest impact on the tungsten reactivity worth.

### Table III

| Reaction channel | Reactivity ρ variation (mk) |
|------------------|----------------------------|
| Radiative capture| -0.28                      |
| Elastic scattering| 0.13                       |
| Inelastic scattering| 0.06                     |
| N,2N             | 0                          |
| N,3N             | 0                          |

According to the above sensitivity analysis results, the reaction rates for the radiative capture, elastic scattering and inelastic scattering in the tungsten target were calculated with FM card in MCNP for case 1 and case 3 in Table II. Integral values of the whole neutron energy interval for the three reaction channel rates are given in the second and third columns of Table IV. From the relative difference RED [RED = (ENDF/B-VII.1 value - ENDF/B-VII.0 value)/ENDF/B-VII.0 value] in the fourth column of Table IV, it can be seen that comparing with case 1, integral value of reaction capture reaction rates in case 3 is larger, whereas that of elastic scattering and inelastic scattering rates in case 3 is smaller. And the relative difference RED of radiation capture reaction rates is maximal among the three reaction channels. Together with the results in Table III, the changes of all three reaction channel rates result in the lower tungsten reactivity value worth in the case 3 and the most significant impact comes from radiation capture reaction channel. Therefore, the cross sections of non-tungsten elements mainly affect the radiative capture, elastic scattering and inelastic scattering rates, which results in the difference of tungsten reactivity worth values from case 1 and case 3. Meanwhile, the difference of radiation capture reaction rates is the main reason for the lower result in case 3.

### Table IV

| Reaction channel | Case 1 | Case 3 | RED (%) |
|------------------|--------|--------|---------|
| Radiative capture| 7.73E-06 | 7.91E-06 | 2.37     |
| Elastic scattering| 5.31E-05 | 5.31E-05 | 0.00    |
| Inelastic scattering| 7.54E-06 | 7.51E-06 | 0.43    |

*Unit: number/source neutron/ cm³

In order to further analyze the influence of non-tungsten elements on main reaction channels, the whole neutrons energy region (10⁸-10 MeV) was divided into 10 energy groups in geometric progression manner of common ratio 10. The relative differences RED of the integral values for the radiative capture, elastic scattering and inelastic scattering reaction rates were respectively calculated in each energy group for case 1 and case 3. As the results in Fig. 3(a-c) shows the maximum RED values of the radiation capture and elastic scattering between the case 1 and case 3 are both presented in the energy interval of 10⁸-10⁵ MeV, whereas the RED of inelastic scattering reaction rates is very low and only found in the energy interval of 10⁵-1 MeV. Therefore, it can be concluded that the difference of non-tungsten elements cross sections between ENDF/B-VII.1 and ENDF/B-VII.0 most mainly results in the discrepancy of the radiation capture and elastic scattering reaction rates in the energy range of 10⁸-10⁵ MeV.
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