The XMM–Newton survey of the Small Magellanic Cloud:
XMMU J005011.2−730026 = SXP 214, a Be/X-ray binary pulsar

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Received 2011 February 18; in original form 2011 January 13

ABSTRACT
In the course of the XMM–Newton survey of the Small Magellanic Cloud, a region to the east of the emission nebula N19 was observed in 2009 November. To search for new candidates for high-mass X-ray binaries, the EPIC-pn and EPIC-MOS data of the detected point sources were investigated and their spectral and temporal characteristics identified. A new transient (XMMU J005011.2−730026 = SXP 214) with a pulse period of 214.05 s was discovered; the source had a hard X-ray spectrum with a power-law index of ≈0.65. The accurate X-ray source location permits the identification of the X-ray source with an ≈15th magnitude Be star, thereby confirming this system as a new Be/X-ray binary.

Key words: stars: neutron – X-rays: binaries.

1 INTRODUCTION AND BACKGROUND
The Be/X-ray systems represent the largest sub-class of all high-mass X-ray binaries (HMXBs). A survey of the literature reveals that of the ≈240 HMXBs known in our Galaxy and the Magellanic Clouds (Liu, van Paradijs & van den Heuvel 2005, 2006), ≥50 per cent fall within this class of binary. In fact, in recent years it has emerged that there is a substantial population of HMXBs in the SMC comparable in number to the Galactic population. Though unlike the Galactic population, all except one of the Small Magellanic Cloud (SMC) HMXBs are Be star systems. In these systems the orbit of the Be star and the compact object, presumably a neutron star, is generally wide and eccentric. X-ray outbursts are normally associated with the passage of the neutron star close to the circumstellar disc (Okazaki & Negueruela 2001) and are generally classified as Type I or II (Stella, White & Rosner 1986). The Type I outbursts occur periodically at the time of the periastron passage of the neutron star, whereas Type II outbursts are much more extensive and occur when the circumstellar material expands to fill most or all of the orbit. This paper concerns itself with Type I outbursts. General reviews of such HMXB systems may be found in Negueruela (1998), Corbet et al. (2009), Coe (2000) and Coe et al. (2009).

One of the aims of the XMM–Newton large programme SMC survey (Haberl & Pietsch 2008a) is the ongoing study of the Be/X-ray binary population of the SMC, which can be used as a star formation tracer for ≈50 Myr old populations (Antoniou et al. 2010). In this paper, we present the analysis of X-ray and optical data from the newly discovered X-ray pulsar XMMU J005011.2−730026, hereafter referred to more simply as SXP 214 following the naming convention of Coe et al. (2005) for X-ray binary pulsars in the SMC.

2 X-RAY OBSERVATIONS
The Local Group galaxies are best suited to study their X-ray source populations using present-day observatories. Extending the existing archival observations, we have carried out a large programme in 2009 with XMM–Newton to obtain a complete X-ray survey of the SMC in the 0.1–10 keV band. The hard X-ray source SXP 214 that is the subject of this work was discovered as transient in observation 14 (Obs ID 0601211401 of the survey), on 2009 November 4. The
field is centred about 22 arcmin east of the emission nebula N19 in the south-western part of the SMC bar. It was observed by all three EPIC instruments (Strüder et al. 2001; Turner et al. 2001) at off-axis angles between 12.9 and 13.7 arcmin on CCDs 10, 3 and 4 for EPIC-pn, EPIC-MOS1 and EPIC-MOS2, respectively. All three cameras were operated in full-frame imaging mode with CCD readout frame times of 73 ms (pn) and 2.6 s (MOS). We used the XMM–Newton Science Analysis System (sas) version 10.0.0\(^1\) to reduce the data.

To correct the astrometric bore-sight, we identified eight sources (mainly known Be/X-ray binaries) in the field of view (FOV) with the Magellanic Clouds Photometric Survey of Zaritsky et al. (2002) and obtained a shift of ΔRA = 0\(^\prime\)10 and ΔDec. = 0\(^\circ\)70. The corrected X-ray position as found by emldetect is RA = 00\(^h\)50\(^m\)1095 and Dec. = -73\(^\circ\)00'25:0 (J2000.0), with a statistical error of 0.18 arcsec and a systematic uncertainty of ∼1 arcsec (1σ confidence for both cases). The relatively large systematic error is caused by the large off-axis angle of the source.

To remove times affected by background flares due to soft protons which occurred at the end of the observation (performed at the end of the satellite orbit), we defined good time intervals by applying thresholds to the background count rate in the 7.0–15.0 keV band of 8 and 2.5 counts ks\(^{-1}\) arcmin\(^{-2}\) for EPIC-pn and EPIC-MOS, respectively. The soft proton background was at a very low level during the first part of the observation, thus resulting in net exposure times of 31.4 and 32.7 ks for EPIC-pn and EPIC-MOS, respectively. The soft proton background was at a very low level during the first part of the observation, thus resulting in net exposure times of 31.4 and 32.7 ks for EPIC-pn and EPIC-MOS, respectively.

To define extraction regions (see Fig. 1) for the source and background with an optimized signal-to-noise ratio, the sas task egression-analyse was used. We ensured that the source extraction region had a distance of > 10 arcsec to other detected sources. For the background extraction, we defined a circular region covering the same point source free area on the sky for all three detectors. To avoid systematic detector background variations present close to CCD borders where the source was located, we chose an area on the same CCDs as the source, and in the case of pn at a similar distance to the readout node. This restricted our selection to a radius of 25 arcsec. We extracted spectra, by selecting single- and double-pixel events from the EPIC-pn data and single-to-quadruple events from EPIC-MOS data, both with LAG = 0. The EPIC-pn/MOS1/MOS2 spectra contained 989/464/396 background-subtracted counts in the 0.2–10.0 keV band, respectively, and were binned to a minimum signal-to-noise ratio of 5 for each bin. To increase the statistics for the timing analysis, we also generated a merged event list from all three instruments, containing 2053 counts (source + background).

2.1 Spectral analysis

We used xspec (Arnaud 1996) version 12.6.0k for spectral fitting. The three EPIC spectra (see Fig. 2) were fitted simultaneously with a common set of spectral model parameters and a relative normalization factor for instrumental differences. The Galactic photoelectric foreground absorption was set to a column density of \(N_{\text{H,GAL}} = 6 × 10^{20}\) cm\(^{-2}\) with abundances according to Wilms, Allen & McCray (2000). The SMC column density was a free parameter with abundances for elements heavier than helium fixed at 0.2 (Russell & Dopita 1992). The emission was modelled with an absorbed power law where we obtained an SMC column density \(N_{\text{H,SMC}} = (4.96 ± 1.50) × 10^{23}\) cm\(^{-2}\); a photon index \(\Gamma = 0.65 ± 0.18\) and a flux in the 0.2–10.0 keV band of \((1.2 ± 0.2) × 10^{-12}\) erg cm\(^{-2}\) s\(^{-1}\). Assuming a distance of 60 kpc, this corresponds to an unabsorbed luminosity of \(7.1 × 10^{35}\) erg s\(^{-1}\). The model resulted in an acceptable fit with \(\chi^2/{\text{dof}} = 57/58\), with relative normalization factors \(c_{\text{MOS1}} = 1.18 ± 0.13\) and \(c_{\text{MOS2}} = 1.03 ± 0.12\) (relative to c\(_{\text{pn}} = 1\)).

Since the EPIC-pn spectrum indicates a weak excess at ∼6.5 keV, we investigated a possible contribution of iron K emission lines at 6.4 keV (fluorescent emission) and 6.7 keV (Fe XXV) by assuming an unresolved linewidth. We obtained 90 per cent upper limits of \(4.0 × 10^{-6}\) and \(3.0 × 10^{-6}\) photons cm\(^{-2}\) s\(^{-1}\), which correspond to equivalent width upper limits of 215 and 163 eV, respectively.

2.2 Timing analysis

The photon arrival times were randomized within the detector CCD frame time and recalculated for the Solar system barycentre using the sas task barycen. We searched for pulsations in the X-ray light

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curves in the EPIC standard energy bands (0.2–0.5, 0.5–1.0, 1.0–2.0, 2.0–4.5 and 4.5–10 keV) and combinations of them, using fast Fourier transform and light-curve folding techniques. The power density spectra derived from light curves in various energy bands from the different EPIC instruments showed a periodic signal at 0.00467 Hz. To maximize the signal-to-noise ratio, we created light curves from the merged EPIC event list (filtered by the GTIs common to EPIC-pn and -MOS). Fig. 3 shows the inferred power density spectrum from the 0.2–10.0 keV energy band with the clear peak at a frequency of 0.00467 Hz. Following Haberl, Eger & Pietsch (2008), we used a Bayesian periodic signal detection method (Gregory & Loredo 1996) to determine the pulse period with a 1σ error to (214.045 ± 0.052) s. In Fig. 4, pulse profiles folded with this period in different energy bands (based on the EPIC standard bands) are shown. Clear variations are seen above 1 keV, while due to the high absorption at lower energies the count rate is insufficient to detect a significant modulation. Hardness ratios were derived from the pulse profiles in two adjacent energy bands \( \frac{R_{i+1} - R_{i}}{R_{i+1} + R_{i}} \) with \( R_i \) denoting the background-subtracted count rate in energy band \( i \) (with \( i \) from 1 to 4). Given the relatively low count rate, no significant hardness ratio variations are seen. We determined a pulsed fraction of \( (29 \pm 9) \% \) for the 0.2–10.0 keV band, assuming a sinusoidal pulse profile.

### 2.3 Long-term X-ray variability

The position of SXP 214 was observed with *XMM–Newton* three times before our large programme SMC survey, without any detection of the source. To derive upper limits, we used sensitivity maps of EPIC-pn and derived 3σ upper limits in the 0.2–10.0 keV band of \( 5.3 \times 10^{-3} \) counts s\(^{-1} \) (Obs ID: 0110000101, 2000 October 15), \( 1.1 \times 10^{-2} \) counts s\(^{-1} \) (Obs ID: 0404680101, 2006 October 5) and \( 4.1 \times 10^{-3} \) counts s\(^{-1} \) (Obs ID: 0403970301, 2007 March 12). Assuming the best-fitting power-law spectrum from above, this corresponds to absorbed flux limits of \( 5.0 \times 10^{-14} \), \( 1.0 \times 10^{-13} \) and \( 3.9 \times 10^{-14} \) erg cm\(^{-2} \) s\(^{-1} \) or absorption-corrected luminosity limits of \( 2.5 \times 10^{34} \), \( 5.0 \times 10^{34} \) and \( 1.9 \times 10^{34} \) erg s\(^{-1} \), respectively. The second limit is higher, due to a shorter background screen exposure of \( \sim 7 \) ks, compared to \( \sim 21 \) ks for the other two observations. The upper limits show that SXP 214 increased in brightness at least by a factor of 30 during its outburst.

SXP 214 has frequently fallen within the FOV of *RXTE* as part of the long-term monitoring of the SMC being carried out by some of the authors on this paper [see Galache et al. (2008) for a discussion of this campaign]. Though the collimator sensitivity to SXP 214 has been very high for periods of many months at a time, there are only eight possible detections with a confidence of \( \geq 99 \) per cent at a high collimator response of \( \geq 0.5 \) in over 10 yr. These detections are shown in Fig. 5. Unfortunately, the sparsity of these detections does not give us any significant insight into the binary period of this system neither do they reveal any obvious long-term spin period change.

The *RXTE* pulse amplitude may be converted to luminosity assuming a distance of 60 kpc to the SMC (though the depth of the sources within the SMC is unknown and may affect this distance by up to \( \pm 10 \) kpc). The X-ray spectrum was assumed to be a power law with a photon index \( \Gamma = 1.5 \) and \( N_H = 6 \times 10^{20} \) cm\(^{-2} \). Furthermore, it was assumed that there was an average pulse fraction of 33 per cent for all the measurements and hence the correct total flux is three times the pulse component. Thus, the luminosity may be determined from the pulse amplitude values shown in Fig. 5 using the relationship

\[
L_X = 0.4 \times 10^{37} \times R \text{ erg s}^{-1},
\]

where \( R \) = pulsed amplitude counts in units of PCU\(^{-1} \) s\(^{-1} \).
The conclusion from the RXTE data is that there have been no major X-ray outbursts ($L_X \geq 10^{37}$ erg s$^{-1}$) over the last decade from SXP 214.

3 OPTICAL COUNTERPART

3.1 Optical and infrared photometry

From its precise X-ray position, the optical counterpart to SXP 214 is identified as SMC SC5 207965 in OGLE II and SMC 100.3 36998 in OGLE III (Udalski, Kubiak & Szymanski 1997) (see Fig. 6). It is also present in the massive compact halo object (MACHO) data as object 212.15966.18, but since these data significantly overlap in time with the OGLE II data, and are not taken through a standard $I$-band filter, they are not discussed any further. The long-term optical measurements of the counterpart to SXP 214 are presented in Fig. 7.

An extra $I$-band point was obtained as part of a set of $BVRI$ observations made using the Faulkes Telescope on 25 November 2009 (MJD 55160), 21 d after the XMM–Newton detection reported here. The Faulkes Telescope is located at Siding Spring, Australia, and is a 2-m, fully autonomous, robotic Ritchey–Chretien reflector on an alt-azimuth mount. The telescope employs a robotic control system. The telescope was used in real-time interface mode for the observation of SXP 214. All the observations were pipeline-processed (flat-fielding and debiasing of the images). The magnitudes of the optical counterpart in all the observed bands were determined by comparison with several other nearby stars on the same image frame. In the case of the $I$-band stars from the OGLE data base, these comparison stars have not exhibited any significant variability in the last 8 yr of OGLE monitoring. In all other wavebands, reference stars from the catalogue of Massey (2002) were used.

The infrared (IR) counterpart is identified as Sirius 00501126−7300260 (Kato et al. 2007) and 2MASS 00501125−7300260. The Sirius observations took place on 31 August 2002 and obtained the following results: $J = 15.32 \pm 0.02$, $K = 15.37 \pm 0.02$. New measurements in the $J$, $H$ and $K$ photometric bands were obtained on 2009 December 11 (MJD 55176) with the same Sirius camera on the 1.4-m IRSF telescope in South Africa (Kato et al. 2007). The full set of photometric measurements is shown in Table 1 and reveal that the IR magnitudes are fainter by $\sim 0.3$ at the time of the XMM–Newton detection compared with the catalogue numbers.
Table 1. Optical and IR photometric measurements of SXP 214 taken in 2009 November and December (see the text for details).

| Waveband | Magnitude | Error on magnitude |
|----------|-----------|---------------------|
| $B$      | 15.22     | 0.02                |
| $V$      | 15.33     | 0.02                |
| $R$      | 15.47     | 0.01                |
| $I$      | 15.47     | 0.03                |
| $J$      | 15.67     | 0.02                |
| $H$      | 15.55     | 0.03                |
| $K$      | 15.61     | 0.09                |

The $(B - V)$ colour index obtained from our data is $(B - V) = -0.11 \pm 0.03$. Correcting for an extinction to the SMC of $E(B - V) = 0.09$ (Schwering & Israel 1991) gives an intrinsic colour of $(B - V) = -0.20 \pm 0.03$. From Wegner (1994), this indicates a spectral type in the range B1V–B3V, typical of optical counterparts to Be/X-ray binaries in the SMC (McBride et al. 2008). However, care must always be taken when interpreting colour information as a spectral type in systems that clearly have circumstellar discs. Such discs can make significant contributions to the $B$ and $V$ bands.

The OGLE III data were detrended using a polynomial function and then subjected to a period search in the range of 2–200 d. Longer periods are harder to explore because they tend to approach the length of each data block ($\sim 200$ d) and the annual observing cycle. Two strong adjacent peaks emerged in the power spectrum (see Fig. 8). These peaks are at 4.520 and 4.576 d, respectively. According to the Corbet diagram for the SMC sources (Corbet et al. 2009), it is very unlikely that SXP 214 could have a binary period of this order. In fact, one of the two peaks (4.576 d) appears to be a beat period between the other period and the annual sampling of the OGLE III data. This is confirmed by examining just 1 yr of the data set which reveals just the one peak at 4.520 d. Furthermore, if this period does not represent a binary period, then the most likely explanation is that it is a non-radial pulsation (NRP) within the Be star (Diago et al. 2008) or, perhaps, the beat of the true NRP period with the daily sampling.

3.2 Optical spectroscopy

Spectroscopic observations of the H$\alpha$ region were made on 11 December 2009 (MJD 55176) using the 1.9-m telescope of the South African Astronomical Observatory (SAAO). A 1200 lines per mm reflection grating blazed at 6800 Å was used with the SITe CCD which is effectively $266 \times 1798$ pixels in size, creating a wavelength coverage from 6200 to 6900 Å. The intrinsic resolution in this mode was 0.42 Å pixel$^{-1}$. The spectrum resulting from a 2000-s integration is shown in Fig. 9.

The signal-to-noise ratio is poor in this spectrum, but there is definite evidence for a weak feature around H$\alpha$. A H$\alpha$ line emission width of $-1.5 \pm 1.0$ Å is estimated. Higher quality data are definitely required to establish if the line shows any structure. Absorption-line features associated with He I at rest wavelengths of 6406 and 6676 Å are also marked in the figure.

Additional observations were made using the integral field spectroscopy (IFS) instrument on 2010 December 28 at Siding Springs Observatory using the 2.3-m Advanced Technology Telescope and its Wide Field Spectrograph (WiFeS). The 1200-s single exposure was made in the central region of SXP 214 at a position angle (east of north) of 0$^\circ$ under clear skies with seeing estimated at 1 arcsec. The exposure was made in classical equal mode using the RT560 beam splitter and 3000 volume phased holographic gratings. For these gratings, the blue (708 lines mm$^{-1}$) range included 3200–5900 Å. The data were reduced using the WiFeS data reduction pipeline based on National Optical Astronomy Observatory IRAF software. This data reduction package was developed from the Gemini IRAF package (McGregor et al. 2003). Use of the pipeline consists of four primary tasks: wifes to set environment parameters; wftable to convert single extension FITS file formats to multi-extension FITS ones and create file lists used by subsequent steps; wfcal to process calibration frames including bias, flat-field, arc and wire; and wfreduce to apply calibration files and create data cubes for analysis. Using QFitsView$^3$, a spectrum of SXP 214 in the blue range was created and is shown in Fig. 10.

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$^3$Written by Thomas Ott and freely available at www.mpe.mpg.de/ott/dpuser/index.html
4 DISCUSSION AND CONCLUSIONS

During the XMM–Newton survey of the SMC, we discovered a new hard X-ray transient with a 0.2–10 keV luminosity of $7 \times 10^{35}$ erg s$^{-1}$. The precise position derived from the EPIC data allowed us to identify a $V = 15.3$ mag early-type star as an optical counterpart. An optical spectrum taken at SAAO revealed a weak H$\beta$ line which clearly identifies the source as a Be/X-ray binary in the SMC. After SXP 11.87, this is the second discovery of such a system during the XMM–Newton SMC survey (Sturm et al. 2011).

SXP 214 shows X-ray pulsations with a period of 214.05 s and a hard (power-law) X-ray spectrum (photon index of $\sim 0.65$). Conspicuous are the residuals in the EPIC-pn spectrum at low energies ($\lesssim 1.5$ keV; see Fig. 2). This might be caused by instrumental differences at high off-axis angles or by the contribution of a soft excess (Hickox, Narayan & Kallman 2004). In the latter case, the low count rate and the high absorption do not allow us to fit an additional component. We note that a soft excess originating near the neutron star (accretion column or disc) and absorbed by the same high column density as the power-law component would require an unrealistic high luminosity of $L \sim 10^{39}$ erg s$^{-1}$ to cause this signal. Therefore, if a soft spectral component exists, it more likely originates in an optically thin gas cloud around the neutron star attenuated by much lower absorption. Given the relatively low X-ray luminosity of SXP 214, this is consistent with the conclusions about the origin of soft excesses in the spectra of HMXBs drawn by Hickox et al. (2004).

At the time of this XMM–Newton detection and other related observations, the source seems to be in an optical low state compared with the past decade. In addition, the three previous XMM–Newton observations of SXP 214 have upper limits at least a factor of 10 below the flux level reported here ($7 \times 10^{35}$ erg s$^{-1}$). So at first glance it seems strange that it was finally detected by XMM–Newton when the optical flux was at such a low point. However, placing SXP 214 on the Corbet diagram would suggest a binary period of the order of 100–200 d. It is therefore very probable that the previous observations failed to catch the system with the neutron star at periastron and hence missed any Type I outbursts that may have occurred. The long-term RXTE monitoring (Fig. 5) at a good collimator sensitivity clearly shows the lack of any more prolonged ($\geq 1$ month) Type II outbursts from this system.

Not only is the $I$-band flux at a very low value for this system, but the H$\alpha$ equivalent width also implies that a very minimal circumstellar disc was present in late 2009. If Kurucz model atmospheres (Kurucz 1979) representing the proposed spectral class range (B1V–B3V) are normalized to the dereddened $B$-band flux [using $E(B-V) = 0.09$ from Schwering & Israel (1991)], it again becomes clear that there is very little evidence for any substantial IR excess arising from a circumstellar disc (see Fig. 11). So anything from B1V with a small disc to B3V with essentially no disc at all offers a satisfactory fit to all the optical and IR data.

A more precise spectral classification may be determined from Fig. 10. Following the method given in Evans et al. (2004), an absence of He I $\lambda\lambda 4026, 4471$ lines defines the class to be later than B0. The weak but approximately equal strength of Mg II ($\lambda 4481$) and Si ii ($\lambda 4553$) narrows it down to B2.5, though it could be up to half a class on either side. In the context of the luminosity class, if we take the $V$-band magnitude presented here and correct for the distance modulus (18.9) and the extinction to the SMC [$E(B-V) = 0.08$], an absolute magnitude of $-3.83$ results. Comparing this to the data in Wegner (1994) predicts a luminosity class III for a B2 star. So the best spectral identification from the data presented here gives B2–B3 III, which adds further support to a minimal circumstellar disc in this system.

In contrast, the column density derived from the X-ray spectrum of $\gtrsim 5 \times 10^{22}$ cm$^{-2}$ is relatively high for Be/X-ray binaries in the SMC (see fig. 12 in Haberl et al. 2008) and largely exceeds the total SMC H$\alpha$ column density of $1.1 \times 10^{22}$ cm$^{-2}$ (Stanimirovic et al. 1999). This suggests a large amount of source-intrinsic absorption close to the neutron star or that the line of sight to the X-ray source during the XMM–Newton observation passed close to the Be star, implying a large system inclination. An even more extreme case is the Be/X-ray binary pulsar SXP 1323, which adds further support to a minimal circumstellar disc in this system.
exceeded $10^{23} \text{ cm}^{-2}$, completely suppressing the hard power-law component at energies below 2 keV and revealing a strong soft spectral component, likely caused by reprocessing in optically thin gas. In the case of SXP 214, this soft component is much weaker, consistent with the above conclusion that only a minimal circumstellar disc was present during the XMM–Newton observation. It is worth noting that within our galaxy columns of up to $10^{24} \text{ cm}^{-2}$ have been measured for some highly obscured HMXB systems.

ACKNOWLEDGMENTS
RS acknowledges support from the BMWI/DLR grant FKZ 50 OR 0907, LJT is supported by a University of Southampton Mayflower Scholarship. As always, we are grateful to the support staff in SAAO for help in using the 1.9-m and IRSF telescopes. The Faulkes Telescope Project is an educational and research arm of the Las Cumbres Observatory Global Telescope Network (LCOGTN). The OGLE project has received funding from the European Community’s Seventh Framework Programme (FP7/2007–2013)/ERC grant agreement no. 246678 to AU. MF thanks Australian Government AINSTO AMNRF for grant number 10/11-O-06.

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