The MAVERIC Survey: A Transitional Millisecond Pulsar Candidate in Terzan 5

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Abstract

Transitional millisecond pulsars are accreting millisecond pulsars that switch between accreting X-ray binary and millisecond radio pulsar states. Only a handful of these objects have been identified so far. Terzan 5 CX1 is a variable hard-X-ray source in the globular cluster Terzan 5. In this paper, we identify a radio counterpart to CX1 in deep Very Large Array radio continuum data. Chandra observations over the last 14 years indicate that CX1 shows two brightness states: in 2003 and 2016 the source was the brightest X-ray source in the cluster (\(L_X \sim 10^{33}\) erg s\(^{-1}\)), while in many intermediate observations, its luminosity was almost an order of magnitude lower. We analyze all available X-ray data of CX1, showing that the two states are consistent with the spectral and variability properties observed for the X-ray active and radio pulsar states of known transitional millisecond pulsars. Finally, we discuss the prospects for the detection of CX1 as a radio pulsar in existing timing data.

Key words: accretion, accretion disks – globular clusters: individual (Terzan 5) – stars: neutron – X-rays: binaries

1. Introduction

1.1. Transitional Millisecond Pulses (MSPs)

MSPs are old radio pulsars that have been spun up through accretion from a companion (Alpar et al. 1982). Neutron star low-mass X-ray binaries (LMXBs) are progenitors of MSPs. This has been confirmed by the detection of coherent millisecond X-ray pulsations in the LMXB SAX J1808.4–3658 (Wijnands \& van der Klis 1998), and later in a dozen other LMXBs (Patruno \& Watts 2012; Patruno et al. 2017). The evolutionary link connecting LMXBs and MSPs has recently been found in the form of transitional millisecond pulsars (tMSPs): neutron stars that move back and forth between accretion-powered LMXB and rotation-powered MSP states on timescales as short as weeks (e.g., Archibald et al. 2009; Papitto et al. 2013). This exciting discovery implies that there is not a simple one-way evolution from LMXBs to MSPs, but instead an extended evolutionary phase in which neutron stars switch between these two different states.

Three tMSPs have now been confirmed through multiwavelength studies; PSR J1023+0038 (Archibald et al. 2009), IGR J18245–2452 in the globular cluster M28 (Papitto et al. 2013), and XSS J12270–4859 (Bassa et al. 2014). All three have shown transitions between a radio pulsar state and a low-level accretion state (e.g., Papitto et al. 2013; Bassa et al. 2014; Stappers et al. 2014; Roy et al. 2015). During the radio pulsar state, radio pulsations are often (but not always) detected, and the radio spectrum is a steep power law (\(S \propto \nu^{-\alpha}\)) with typical spectral indexes similar to those of other MSPs, \(\alpha \sim -1.6\) to \(-1.8\) (e.g., Kramer et al. 1998), while the X-ray emission is relatively faint (\(L_X \sim 10^{32}\) erg s\(^{-1}\)), Bogdanov et al. 2014; Linares 2014; Roberts et al. 2015). In the low-level accretion state, radio pulsations cease, and bright radio emission with a flat radio spectrum (\(\alpha \sim 0\)) and rapid variability (e.g., Deller et al. 2015) are thought to be produced by a partially self-absorbed synchrotron jet, while the X-ray luminosity increases to \(L_X \sim 10^{33–34}\) erg s\(^{-1}\) (e.g., Bogdanov et al. 2015). During this state, X-ray pulsations thought to be powered through accretion have been detected (Papitto et al. 2013, 2015; Archibald et al. 2015; Jaodand et al. 2016). In both states the X-ray spectrum is relatively hard, with power-law photon index \(1–1.5\) (\(\frac{\mathrm{d}N}{\mathrm{d}E} \propto E^{-1}\)). Furthermore, in the low-level accretion state, tMSPs tend to show significant variability, and occasional flares (e.g., Takata et al. 2014; Tendulkar et al. 2014; Papitto et al. 2015). They can also show two dominant “modes” of brightness (labeled active and passive modes), whose flux can differ by a factor of a few (e.g., Bogdanov et al. 2014).

It is worth noting that radio emission from tMSPs during the low-level accretion state is thought to be due to a partially self-absorbed jet (Deller et al. 2015). However, Bogdanov et al. (2018) recently showed a sharp anti-correlation between radio and X-ray brightness in PSR J1023+0038 during this state. This is at odds with the standard accretion-jet coupling. The synchronous rise and decay of the radio emission during the drop into, and rise out of, the X-ray low mode, suggests that the radio-emitting region is causally coupled to the jets, which implies it must be very compact (tens of light seconds). Furthermore, Ambrosino et al. (2017) recently detected optical pulsations from PSR J1023+0038 during the low-level accretion state, which indicates that the optical (and perhaps X-ray) emission in this state from these systems may be powered through magnetospheric processes.

Correlated changes in the X-rays and radio, along with changes in the optical and gamma-rays (harder to study in distant, crowded globular clusters) have clearly identified X-ray high and low transitions in the nearby (1–2 kpc) tMSPs PSR J1023+0038 (Patruno et al. 2014; Stappers et al. 2014) and
XSS J12270–4859 (Bassa et al. 2014; Roy et al. 2015). Similar tMSP transitions have not yet been used to identify tMSPs at greater distances, where radio pulsar searches are more difficult; the M28 tMSP was identified when it showed a uniquely bright ($\sim 10^{36}$ erg s$^{-1}$) accretion-powered X-ray outburst, and then returned to a radio pulsar state (Papitto et al. 2013).

In addition to the three confirmed tMSPs discussed above, there are a handful of candidate tMSPs discovered through multi-wavelength studies. These tMSP candidates include 1FGL J1417.7–4407, a gamma-ray source with an MSP counterpart and spectroscopic signs of mass transfer (Strader et al. 2015; Camilo et al. 2016); 1RXS J154439.4–112820 (3FGL J1544.6–1125), a gamma-ray and X-ray binary showing similar X-ray mode-switching behavior as seen in confirmed tMSPs (Bogdanov & Halpern 2015; Bogdanov et al. 2016; Britt et al. 2017); and 3FGL J0427.9–6704, an eclipsing low-mass X-ray and gamma-ray binary with rapid X-ray variability (Strader et al. 2016).

Redbacks are MSPs in which the companion is a hydrogen-rich star showing evidence of mass loss, with mass $\geq 0.2–0.5 M_\odot$ (Roberts 2013). In comparison, black widows are MSPs with mass-losing companions with mass $\ll 0.1 M_\odot$ (Fruchter et al. 1988; Roberts 2013). All three of the confirmed tMSPs so far are redback systems (Archibald et al. 2013; Papitto et al. 2013; Roy et al. 2015). This has motivated a search for tMSP systems among redbacks and black widow pulsars. Numerous black widows and redbacks have high-energy counterparts in X-rays and gamma-rays (e.g., Ray et al. 2012; Roberts 2013; Gentile et al. 2014; Linares 2014; Roberts et al. 2015). In X-rays, redbacks are typically brighter ($L_\gamma \sim 10^{31–32}$ erg s$^{-1}$) and show a harder non-thermal X-ray spectrum in comparison with black widows, where the nonthermal X-rays are fainter ($L_\gamma \sim 10^{29–31}$ erg s$^{-1}$) and sometimes dominated by thermal X-rays from the MSP surface (Gentile et al. 2014; Roberts et al. 2015). It is thought that the nonthermal X-rays in these systems mostly come from an intrabinary shock between the pulsar wind and material flowing from the companion (Bogdanov et al. 2005; Romani & Sanchez 2016).

### 1.2. Terzan 5

Terzan 5 is a massive ($\sim 10^6 M_\odot$) globular cluster located $\sim 5.9$ kpc away in the Galactic bulge (Valenti et al. 2007; Lanzoni et al. 2010). It contains one of the largest populations of X-ray sources ($\geq 50$; Heinke et al. 2006b) and the highest number of radio MSPs (Prager et al. 2017) of any Galactic globular cluster. This high population of X-ray and MSP binaries in Terzan 5, compared to other globular clusters, is attributed to dynamics. The formation rate of X-ray binaries per unit mass in globular clusters is roughly two orders of magnitude higher than that in the Galactic field (Katz 1975). It has been shown that this is due to the higher stellar density in these clusters, which leads to the formation of X-ray binaries through dynamical encounters (Verbunt & Hut 1987; Heinke et al. 2003b; Pooley et al. 2003; Jordán et al. 2004; Bahramian et al. 2013), as opposed to formation through isolated binary evolution, which is thought to be the dominant channel in the Galactic field.

Terzan 5 CX1 (CXOGLb J174804.5–244641) is a variable hard X-ray source in the core of Terzan 5. It was first identified in Chandra observations in 2003, when it was the brightest source in the cluster (Heinke et al. 2006b). The nature of this source is not yet fully understood. In this paper, we present analyses of all the available Chandra data on this source. We also identify a radio continuum counterpart, using data taken as part of the MAVERIC (Milky-way ATCA and VLA Exploration of Radio-sources in Clusters; Tremou et al. 2018; L. Shishkovsky et al. 2018, in preparation; V. Tudor et al. 2018, in preparation) survey. Based on the X-ray variability, hard-X-ray spectrum and the possibly steep spectrum radio counterpart, we suggest a tMSP nature for this object.

We present the data, reduction methods, and analysis in Section 2, and our results in Section 3. Finally, we discuss possible scenarios for the nature of the source in Section 4. In our analysis we assume a distance of 5.9 kpc to Terzan 5 (Valenti et al. 2007). All uncertainties reported/plotted based on observations presented in this work have $\pm 68\%$ confidence.

### 2. Data Reduction and Analysis

#### 2.1. X-Rays: Chandra

Terzan 5 has been observed with Chandra numerous times since 2001, as it contains a large number of X-ray binaries, three of which are transients that have shown outbursts with $L_\gamma > 10^{36}$ erg s$^{-1}$ (Heinke et al. 2003a; Wijnands et al. 2005; Strohmayer & Markwardt 2010; Bahramian et al. 2014). We reduced and analyzed all archival Chandra/ACIS observations of Terzan 5 in which all cluster sources are faint/quiet (Table 1). This excludes observations taken during the outbursts of transient XRBs within Terzan 5 (Obs. IDs 664, 655, 11051, 12454, 13708). Additionally, we analyzed observations from our 2016 campaign (PI: N. Degenaar). All Chandra observations in this study were performed with ACIS-S in faint mode, with Terzan 5 on chip S3, and the cluster core on-axis.

| Obs. ID | Date        | Exposure (ks) | Avg. Rate (cts/s) |
|---------|-------------|---------------|------------------|
| 03798   | 2003 Jul 13 | 39            | $1.17(\pm0.05) \times 10^{-2}$ |
| 10059   | 2009 Jul 15 | 36            | $5.9(\pm1.4) \times 10^{-4}$  |
| 13225   | 2011 Feb 17 | 30            | $4.2(\pm1.4) \times 10^{-4}$  |
| 13252   | 2011 Apr 29 | 40            | $9.6(\pm1.6) \times 10^{-4}$  |
| 13705   | 2011 Sep 05 | 14            | $8.6(\pm2.7) \times 10^{-4}$  |
| 14339   | 2011 Sep 08 | 34            | $9.3(\pm1.8) \times 10^{-4}$  |
| 13706   | 2012 May 14 | 46            | $6.8(\pm1.3) \times 10^{-4}$  |
| 14475   | 2012 Sep 17 | 30            | $2.9(\pm1.2) \times 10^{-4}$  |
| 14476   | 2012 Oct 28 | 29            | $3.0(\pm1.2) \times 10^{-4}$  |
| 14477   | 2013 Feb 05 | 29            | $5.3(\pm1.5) \times 10^{-4}$  |
| 14478   | 2013 Jul 17 | 29            | $4.4(\pm1.4) \times 10^{-4}$  |
| 14625   | 2013 Feb 22 | 49            | $6.5(\pm1.2) \times 10^{-4}$  |
| 15615   | 2013 Feb 23 | 84            | $4.5(\pm0.8) \times 10^{-4}$  |
| 14479   | 2014 Jul 15 | 29            | $5.7(\pm1.7) \times 10^{-4}$  |
| 16638   | 2014 Jul 17 | 72            | $6.3(\pm1.0) \times 10^{-4}$  |
| 15750   | 2014 Jul 20 | 23            | $7.8(\pm1.9) \times 10^{-4}$  |
| 17779   | 2016 Jul 14 | 69            | $8.8(\pm0.4) \times 10^{-3}$  |
| 18881   | 2016 Jul 15 | 65            | $8.2(\pm0.4) \times 10^{-3}$  |

Note. Net count rates are in the 0.5–10 keV band. We omitted all observations taken during the outburst of transient XRBs within Terzan 5 (Obs. IDs 664, 655, 11051, 12454, 13708).
We used CIAO 4.9 and CALDB 4.7.3 (Fruscione et al. 2006) for data reprocessing and analysis. We reprocessed all Chandra data with chandra_repro and extracted source and background spectra using specextract. CX1 is located in the crowded core of Terzan 5, ∼5″ from the center of the cluster (Figure 1). Thus, we used a circular extraction region for the source with a radius of 1″ to avoid contamination from nearby sources, and an annulus with inner and outer radii of 1.5 and 6″ for the background, while excluding all detected sources in this region. Finally, we used XSPEC 12.9.1n (Arnaud 1996) for spectral analysis.

2.2. Radio Continuum: VLA

Deep radio continuum observations of Terzan 5 were obtained with the NSF’s Karl G. Jansky Very Large Array (VLA) in 2012 and 2014. The 2012 observations were taken as part of the MAVERIC survey to study radio sources in globular clusters (L. Shishkovsky et al. 2018, in preparation; V. Tudor et al. 2018, in preparation), while the 2014 data were obtained to study a Terzan 5 pulsar. In addition to these previously unpublished data, we also use shallower VLA observations from 2015 obtained to study an outburst of the Terzan 5 low-mass X-ray binary EXO 1745–248 (Tetarenko et al. 2016).

For the 2012 observations we observed for five hours at C-band on September 10–12 (with two 1 GHz basebands centered at 5.0 and 7.5 GHz) and five hours at S-band on September 13–14 (with the two 1 GHz basebands centered at 2.5 and 3.5 GHz). The VLA was in BnA configuration. This bandwidth was sampled with 2 MHz wide channels and was obtained in full Stokes mode. Calibration was carried out using the complex gain calibrator J1751–2524, while 3C48 provided absolute flux calibration and bandpass calibration.

On 2014 February 14, 3.5 hr of S-band data were obtained on Terzan 5 while the VLA was in BnA-to-A move configuration. These data were observed under test program TPUL0001 (PI P. Demorest). The bandwidth was 1024 MHz centered at 2.97 GHz, sampled by 0.5 MHz wide channels and with dual polarization. The complex gain calibrator was J1751–2524, while 3C48 provided absolute flux calibration and bandpass calibration.

Data editing and calibration were carried out using standard routines in AIPS (Greisen 2003). For imaging, the bandwidth was split in half in order to minimize frequency-dependent imaging artifacts. Imaging was performed on each frequency chunk separately, using the task IMAGR in AIPS with a Briggs robust value of 1. We self-calibrated the 2012 data in order to minimize artifacts around bright sources. Table 2 lists the
imaging parameters achieved. The 2012 C-band data were observed at low elevation, so the synthesized beam is rather elongated.

The reduction and analysis of the 2015 VLA data is described in Tetarenko et al. (2016). These data comprise three 45-minute X-band blocks observed in B configuration, from 2015 March 19 to April 12. Each block had two 2048 MHz subbands centered at 9.0 and 11.0 GHz. The results presented here use all the data for each subband imaged together.

3. Results

3.1. X-Ray Variability and Spectral Analysis

Comparing Chandra observations of CX1, the source shows clear variations in brightness between observations (Table 1, Figure 2). While in 2003 and 2016 CX1 has an X-ray luminosity of \( L_X \gtrsim 10^{35} \text{ erg s}^{-1} \), in all observations between 2009 and 2014 it shows an average X-ray luminosity of \( \sim 1.3 \times 10^{32} \text{ erg s}^{-1} \). This indicates that CX1 shows distinct bright and faint “states.” As we discuss in Section 4, these states resemble those seen in tMSP systems.

Besides changes in brightness on month/year timescales, CX1 also shows variability by a factor of a few within observations on timescales of minutes/hours, clearly visible in the 2003 and 2016 observations, when the source was in the bright state (Figure 3).

There is no evidence for variability in the observations taken during the faint state, using a \( \chi^2 \) test. However, this apparent lack of variability could be due to our low signal-to-noise ratio in the faint-state observations, caused by the large distance and extinction toward Terzan 5 (Terzan 5 is just slightly further away than M28, the host cluster of the tMSP IGR J18245–2452, but it suffers almost 10 times more extinction).

We test this possibility by artificially degrading CX1’s light curves from observations taken during the bright state (2003 and 2016) so that the average source count rate in these observations match the average rate estimated in the faint state. This degradation increases the uncertainties estimated on counts (and subsequently rates) significantly, enough to make the variations undetectable in a \( \chi^2 \) test.

We used XSPEC 12.9.1 (Arnaud 1996) for spectral analysis. We assumed Wilms et al. (2000) elemental abundances and Verner et al. (1996) photoelectric absorption cross sections.

While spectra taken during the bright state (in 2003 and 2016) each have more than 400 counts, the spectra from observations during the faint state (2009–2014) have only a few counts (\( \sim 10 \) to 50). Since we have no evidence of variation between these observations (Table 1, Figure 2), we combined all spectra from these observations using the CIAO task COMBINE_SPEC. We binned all the final spectra (3 during the bright states in 2003 and 2016, and the combined spectrum of observations between 2009 and 2014) by 20 counts per bin, and used chi-squared statistics for fitting and goodness-of-fit tests.

We fit each of these four spectra separately with an absorbed power law (tbabs x pegpwrlw in XSPEC). As we find no evidence for intrinsic or variable absorption in this system, we set the hydrogen column density \( (N_H) \) to the value estimated by X-ray spectral fitting for the cluster, \( 2.6 \times 10^{22} \text{ cm}^{-2} \) (Bahramian et al. 2015).

The results of our spectral analysis are tabulated in Table 3 and the spectra are plotted in Figure 4. A power-law model fits all the spectra well. While all the spectra are hard, there is clear evidence for spectral variability. The spectrum seems hardest in the faint state, with a photon index of \( 0.7 \pm 0.2 \). It is worth noting

### Table 2

| Date          | Freq (GHz) | Major Axis (arcsec) | Minor Axis (arcsec) | Image rms Noise (\( \mu \text{Jy/beam} \)) | Source Flux (\( \mu \text{Jy} \)) |
|---------------|------------|---------------------|---------------------|------------------------------------------|-------------------------------|
| 2012 Sep 10–14| 2.6        | 2.35                | 2.05                | 5.7                                      | …                            |
| 2012 Sep 10–14| 3.4        | 1.64                | 1.39                | 5.5                                      | <10.6                        |
| 2012 Sep 10–14| 5.0        | 2.48                | 0.79                | 4.7                                      | 16.4 ± 4.9                   |
| 2012 Sep 10–14| 7.4        | 1.73                | 0.47                | 3.6                                      | <10.6                        |
| 2014 Feb 14   | 2.7        | 1.18                | 0.59                | 5.0                                      | 17.7 ± 5.7                   |
| 2014 Feb 14   | 3.2        | 1.00                | 0.50                | 5.5                                      | 22.0 ± 5.9                   |
| 2015 Mar/Apr  | 9.0        | 1.60                | 0.80                | 5.0                                      | <14.9                        |
| 2015 Mar/Apr  | 11.0       | 1.28                | 0.66                | 5.8                                      | <17.5                        |

Note. Major axis and minor axis are FWHM values for the synthesized beam. We were not able to constrain the source flux in S band data from 2012, as the source appears blended with the bright nearby pulsar Terzan 5 M. Source flux upper limits in this table are 3\( \sigma \). The 2015 observations were taken in portions on March 19th and 24th and April 12th.
To average count rates in the relevant observation. To convert count rates to luminosities, data points in each light curve were multiplied by the ratio of unabsorbed flux (as obtained from spectroscopy, see Section 3.1) to average count rate in the relevant observation.

Finally, we note that a precise X-ray position of CX1, tied to a high-quality astrometric frame, is required for multi-wavelength matching of possible counterparts. We start from the coordinates given in the X-ray catalog of Heinke et al. (2006b). To check the absolute astrometry, we note that there are clear X-ray matches with two known MSPs with accurate positions from radio timing (Prager et al. 2017). In addition to the match between CX10 and the pulsar Terzan 5 P noted in Heinke et al. (2006b), the source CX13 appears to be the counterpart of pulsar Terzan 5 ad (analysis of the X-ray counterparts of Terzan 5 MSPs will be published separately; S. Bogdanov et al. 2018, in preparation). These two matches show evidence for a consistent offset of about $-0^\circ12$ in R.A. and $+0^\circ29$ in decl. When these offsets are applied to the CX1 position listed in Heinke et al. (2006b), we find an updated J2000 position of (R.A., decl.) = (17:48:04.578, $-24:46:42.24$), with an uncertainty of $0^\circ15$, which we use for the remainder of the paper. A deeper catalog of X-ray sources in Terzan 5, useful for finding additional matches, is beyond the scope of this paper (see A. Bahramian et al. 2018, in preparation).

3.2. Radio Counterpart: Spectrum and Variability

CX1 is clearly detected as a radio source in a subset of the VLA observations, but is marginally or not detected in others. The 2014 S-band data were taken in the most extended configuration of our VLA data sets and have the best angular resolution. CX1 is detected in both subbands, with flux densities of $17.7 \pm 5.7 \mu$Jy (2.7 GHz) and $22.0 \pm 5.9 \mu$Jy (3.2 GHz). At 2.7 GHz the J2000 position of this radio source is (R.A., decl.) = (17:48:04.575 ± 0\"08, $-24:46:42.12$ ± 0\"16), matching the X-ray position of CX1 to within 0\"13. At 3.2 GHz the radio source is at (R.A., decl.) = (17:48:04.587 ± 0\"06, $-24:46:42.28$ ± 0\"12), also about 0\"13 from the position of CX1, even though the 3.2 GHz position is possibly affected by a faint source located southeast of the main source (this faint source is not detected at 2.7 GHz). In any case, the excellent astrometric match between CX1 and the radio source in both subbands (separation of less than 0\"15, or 1\sigma, see Figure 1) strongly suggests it is the counterpart to CX1.

| Data Set | $\Gamma$ | Flux | $\chi^2$/d.o.f | N.H.P. |
|----------|----------|------|---------------|--------|
| 2003     | $1.1 \pm 0.1$ | $6.1(\pm0.4) \times 10^{-13}$ | 0.55/20 | 94.4    |
| 2009–2014| $0.7^{+0.2}_{-0.1}$ | $3.3(\pm0.3) \times 10^{-14}$ | 1.04/18 | 41.0    |
| 2016A    | $1.5 \pm 0.1$ | $4.6(\pm0.2) \times 10^{-13}$ | 0.85/27 | 70.9    |
| 2016B    | $1.4 \pm 0.1$ | $4.2(\pm0.2) \times 10^{-13}$ | 1.09/24 | 34.9    |

Table 3
Results of X-Ray Spectral Analysis

Note. All uncertainties have 68% confidence. Unabsorbed flux reported in the 0.5–10 keV band in erg s$^{-1}$ cm$^{-2}$. N.H.P. is XSpec’s null hypothesis probability in percent. Data set 2009–2014 is the combined spectrum from all observations between 2009 and 2014 (see Section 2.1). The 2016A spectrum is from an observation on 2016 July 14 (obs.ID 17779) and 2016B is from 2016 July 15 (obs.ID 18881).

that the spectra from the bright states in 2003 and 2016 are not completely similar either. In 2003, the spectrum appears comparatively harder (with a photon index of $1.1 \pm 0.1$), and the source seems slightly brighter.

Finally, we note that a precise X-ray position of CX1, tied to a high-quality astrometric frame, is required for multi-wavelength matching of possible counterparts. We start from the coordinates given in the X-ray catalog of Heinke et al. (2006b). To check the absolute astrometry, we note that there are clear X-ray matches with two known MSPs with accurate positions from radio timing (Prager et al. 2017). In addition to the match between CX10 and the pulsar Terzan 5 P noted in Heinke et al. (2006b), the source CX13 appears to be the counterpart of pulsar Terzan 5 ad (analysis of the X-ray counterparts of Terzan 5 MSPs will be published separately; S. Bogdanov et al. 2018, in preparation). These two matches show evidence for a consistent offset of about $-0^\circ12$ in R.A. and $+0^\circ29$ in decl. When these offsets are applied to the CX1 position listed in Heinke et al. (2006b), we find an updated J2000 position of (R.A., decl.) = (17:48:04.578, $-24:46:42.24$), with an uncertainty of $0^\circ15$, which we use for the remainder of the paper. A deeper catalog of X-ray sources in Terzan 5, useful for finding additional matches, is beyond the scope of this paper (see A. Bahramian et al. 2018, in preparation).

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In the 2012 S-band VLA data (with subbands centered at 2.6 and 3.4 GHz after RFI flagging) there is marginal evidence of a source at the position of CX1, but it is blended with the bright nearby pulsar Terzan 5 M, and we were unable to obtain reliable flux densities, even when simultaneously fitting for the flux density of the pulsar.

In the 2012 C-band data, we detected a radio source at the position of CX1 at 5.0 GHz. Its flux density is $16.4 \pm 4.9 \, \mu$Jy. The source was not detected at 7.4 GHz, with a $3\sigma$ upper limit of $10.6 \, \mu$Jy.

None of the individual-epoch radio observations alone provide meaningful constraints on the radio spectral index of the source. Under the assumption that the source did not vary during the time period of the VLA observations (all of which were in the X-ray faint state), we can constrain the spectral index by modeling the individual detections as well as the upper limits. For a power law with $S \propto \nu^{\alpha}$ (where $S$ is the flux density, $\nu$ is the frequency, and $\alpha$ is the spectral index), we find $\alpha = -1.05^{+0.48}_{-0.65}$, where the uncertainties represent a 68% confidence interval. This spectral index does not strongly constrain the nature of the source: it is consistent within 1$\sigma$ with the steep negative spectral indices typical of MSPs ($\alpha \sim -1.6$ to $-1.8$; Kramer et al. 1998), but is also consistent within 2$\sigma$ with the flat spectra expected for self-absorbed compact jets ($\alpha \sim 0.0$–0.6).

One concern in these results is the potential contamination from the nearby source in the 2014 data (see Figure 1). The peak flux density for CX1 from the image is consistent with the fitted value within the uncertainty, so there is no evidence that the flux density measurement is significantly biased. To the extent that it has an effect, it would probably bias the flux density slightly upward, which would tend to suggest a flatter spectral index. To quantify this, we took the peak flux density of the central pixel of CX1 (18.5 $\mu$Jy) and re-ran the spectral index analysis. We find a value of $-1.02^{+0.51}_{-0.71}$, identical (within uncertainties) to the value mentioned above. Hence, there is no evidence that this issue has an effect on any of our conclusions.

3.3. Constraints on an Optical Counterpart

Terzan 5 was observed by HST/ACS in 2003, 2013, and 2015 with the F606W and F814W filters. We reduced and analyzed these images from the Hubble Legacy Archive\(^9\) to search for an optical counterpart.

We set the absolute astrometry of the HST images by performing source detection on every image using the daofind task provided in PYRAF 2.16 (Science Software Branch at STScI 2012) and cross-matching to the Gaia data release 1 source catalog (Gaia Collaboration et al. 2016) using the task ccmpair. This provides mean rms values of 0.03 arcsec, 0.009 arcsec, and 0.007 arcsec for 2003, 2013, and 2015 epochs respectively, which we consider representative of the absolute astrometric uncertainty.

The density of stars in the core of Terzan 5 is extremely high, and there is a source at R.A. = 17:48:04.575, decl. = $-24:46:42.310$, formally $0\partial08$ from CX1 (Figure 5). Aperture photometry of this source yields mean $m_{F606W}$ and $m_{F814W}$ magnitudes of $24.4 \pm 0.5$ and $21.3 \pm 0.3$ mag (Vega mag), respectively. There is no significant evidence of variation across the HST epochs. Compared to the published color–magnitude diagrams of Terzan 5 in these filters (e.g., see Massari et al. 2012), this photometry is consistent with the identification of this star as a likely main sequence star, though the high and variable extinction makes a more precise classification impossible. The optical (and especially infrared) counterparts of many X-ray binaries in quiescence appear quite like “normal” stars, as the optical/infrared emission is completely dominated by the donor star. However, at least two of the HST epochs targeting Terzan 5 are taken around times when CX1 shows brighter (few $10^{33}$ erg s$^{-1}$) X-ray emission (in 2003 and 2015). Transitional MSPs typically show increases in brightness of 1–2 when they change to the bright state (Halpern et al. 2013; Bassa et al. 2014). Thus, the “normal” color of this star, as well as its lack of variability over the observed epochs (which are taken around times the source is bright in X-rays), suggests that it is unlikely to be the true counterpart to CX1.

To estimate the probability of a random match with an optical source in the core of Terzan 5, we picked 50,000 random coordinates within the core of Terzan 5 and measured the distance to the closest real detection from each (in the deepest HST image). The distributions of these separations appear Gaussian, with a peak around $0.018$ and a standard deviation of $0.009$. This distribution indicates that the probability of finding an optical source $<0.15$ (uncertainty radius in the X-ray coordinates) is $\sim35\%$. This reinforces the conclusion above that the putative optical counterpart might well be an interloper and that the true counterpart may be too faint to be detected in these data.

4. Discussion

4.1. Terzan 5 CX1 as a tMSP

In X-rays, Terzan 5 CX1 shows two brightness states. The faint state appears to last for years (considering the many consecutive faint-state observations), and the few observations of the bright state suggest a duty cycle of $\sim1/6$–$1/3$. In the bright state, the source shows an average luminosity of $L_x \approx 2 \times 10^{33}$ erg s$^{-1}$, and in the faint state a luminosity of $L_x \approx 1.3 \times 10^{32}$ erg s$^{-1}$ in the 0.5–10 keV band. These luminosity levels and their timescales are consistent with the low-level accretion (bright) and pulsar (faint) states observed in the tMSP systems. During the pulsar state, PSR J1023+0038 shows an X-ray luminosity between $9 \times 10^{31}$ and $1.6 \times 10^{32}$ erg s$^{-1}$ in the 0.5–10 keV band (Archibald et al. 2010; Bogdanov et al. 2011; Linares 2014). In the same state, XSS J12270–4859 shows $L_x \approx 1.6 \times 10^{32}$ (Bogdanov et al. 2014; Papitto et al. 2015) and IGR J18245–2452 shows $L_x \approx 2.2 \times 10^{32}$ (Linares et al. 2014). These luminosities are typical of redback rotation-powered pulsars (Gentile et al. 2014; Linares 2014; Roberts et al. 2015).

In the low-level accretion state, tMSPs tend to show significant variability, and occasional flares (e.g., Takata et al. 2014; Tendulkar et al. 2014; Papitto et al. 2015). They can also show two dominant “modes” of brightness (labeled active and passive modes), which can be different by a factor of a few. These modes can last hours and are distinct enough that the distribution of X-ray brightness appears as a bimodal distribution (e.g., De Martino et al. 2013; Linares et al. 2014; Bogdanov et al. 2015).

Terzan 5 CX1 shows clear variability during the bright state. A KS test indicates a probability of constancy of $10^{-16}$ for observations during the bright state. Some of these variations

\(^9\) http://hla.stsci.edu/
appear to be flares and at times the system stays at one brightness level for a few ks (Figure 3), not unlike modes observed in tMSPs. However, the X-ray brightness does not show a clear bimodal distribution (Figure 6, left). However, there is a tentative peak in the histogram at \( \log(L_X) \approx 32.5 \) erg s\(^{-1}\), hinting at a possible bimodal distribution. The absence of a clear bimodal flux distribution in this case could be due to sensitivity, as for PSR J1023+0038—which is significantly closer and less extincted than Terzan 5—Bogdanov et al. (2015) found the longest low mode to be 1.9 ks, and the longest high mode 6.9 ks, while most of the periods in a single mode are well under 1 ks (see their Figure 6). This is shorter than the timescales we are sensitive to in Chandra observations of Terzan 5; see Figure 3.

To test this further, we downloaded the PSR J1023+0038 XMM-Newton/EPIC light curve from obsID 0720030101, which shows numerous low and high modes (reported in Bogdanov et al. 2015), and rebinned the light curve to 1 ks per bin (chosen as a timescale similar to our sensitivity levels in the Chandra light curve of Terzan 5 CX1). Doing so completely removes the bimodal distribution. However, while PSR J1023+0038 and XSS J12270−4859 tend to show short low modes, the low mode has been observed to last for \( \sim 10 \) ks in IGR J18245−2452 (Linares et al. 2014). Thus, we recognize that a key aspect of the X-ray phenomenology of known tMSPs, a bimodal count rate in

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**Figure 5.** HST/ACS images of Terzan 5 CX1 (red circle). The source in CX1’s error circle (indicated with a green diamond) appears to be a “normal” main-sequence star and is unlikely to be the counterpart (see Section 3.3 for details).

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**Figure 6.** Left: distribution of brightness in Terzan 5 CX1. The blue data set represents the distribution based on observations during the bright state (similar to the low-level accretion state in tMSPs), while the pink data set shows the distribution during the faint state (similar to the pulsar state in tMSPs). There is no clear bimodality visible in the bright state indicating the presence of high and low modes, as have been seen in tMSPs. Luminosity is in the 0.5–10 keV band in erg s\(^{-1}\). Right: hardness–luminosity diagram of redbacks (yellow triangles), black widows (purple pentagons), and low-level accretion states in tMSPs with modes separated (active mode represented by blue circle and passive mode by pink squares). The dashed lines show the evolution of tMSPs between the accretion state and the pulsar state. Terzan 5 CX1 is shown by orange diamonds. In the bright state, Terzan 5 CX1 behavior is consistent with average tMSP luminosity and hardness in the accretion state. In the faint state, it is very similar (slightly harder) to redbacks. The plot is adapted from Linares (2014). Data for PSR J2129-0429 are from Roberts et al. (2015). Data for black widows and PSR J2215+5135 are from Gentile et al. (2014).
the disk state could only be detected in Ter5 CX1 if it follows the longer-duration pattern seen in IGR J18245–2452.

tMSPs (in the low-level accretion state and the pulsar state) and redback pulsars tend to show hard X-ray spectra, and are generally well-described by a power-law model with photon index $\leq 1.6$ (e.g., Linares 2014; Roberts et al. 2015). On average, the X-ray spectra of tMSPs are strikingly similar between the different states (as opposed to NS LMXBs, for instance, which can change significantly in slope over 1 mag in X-ray luminosity; see, e.g., Wijnands et al. 2015). Terzan 5 CX1’s spectra seem to follow a similar trend as the tMSPs (Figure 6, right). While the bright state in this system cannot be clearly broken into “high” and “low” modes, the hardness and luminosity are consistent with the average luminosity and hardness of tMSPs in the low-level accretion state, and the luminosity and hardness of the faint state of this system are consistent with those of redbacks (including tMSPs).

The 2012 VLA data were taken within a week of a Chandra observation finding CX1 in its faint state, strongly indicating it was X-ray faint during those observations. The 2014 VLA data were 5 months from the nearest Chandra observation, but as all Chandra observations within 2 years of this observation found CX1 to be X-ray faint, we reasonably assume these VLA data were also taken in the faint state. In the tMSP interpretation, the X-ray faint state implies a radio pulsar state, which is indeed slightly preferred by our limited constraints on the radio spectral index. The 2015 VLA data were observed between a 2014 Chandra observation in the X-ray faint state, and a 2016 Chandra observation in the bright state, so we do not know CX1’s X-ray state in 2015. The upper limits from the 2015 VLA data are not constraining for either hypothesis. Assuming that, in the bright state, CX1 has a radio-X-ray luminosity comparable to that of the tMSP PSR J1023 + 0023 (Deller et al. 2015), then the predicted mean radio flux density in the 2015 VLA observations would be $\sim 5 \mu Jy$, too faint to be detected in the 2015 VLA data (which had a stacked rms noise of $3.8 \mu Jy/beam$).

### 4.2. Other Possibilities for the Nature of Terzan 5 CX1

Terzan 5 CX1 is located in the core of Terzan 5, $\sim 5^\prime$ from the center of the cluster. This makes the likelihood of a chance alignment extremely small: it is almost certainly associated with the cluster. This dramatically reduces the possibility of it being an AGN or a foreground flare star. The absence of a bright optical counterpart, and the absorption of its X-ray spectrum, also argue against a foreground flare star nature.

CX1’s X-ray luminosity in the bright state is also similar to that of the brightest cataclysmic variables (CVs, e.g., Stacey et al. 2011; Bernardini et al. 2012; Pretorius & Mukai 2014). Many CVs show detectable variability in X-rays and a hard X-ray spectrum. However, in the radio, CX1’s persistent emission is brighter than that of the brightest CVs. An example is SS Cyg, a CV with similar X-ray luminosity as CX1 during the faint state, which shows bright radio flares up to $\sim 1.4 \times 10^{27}$ erg s$^{-1}$ at 5 GHz (Russell et al. 2016; Mooley et al. 2017). While these flares show radio luminosities similar to that of CX1, they are very short-lived (lasting $\sim$minutes), while CX1 appears persistent, as measured in multiple epochs.

Many CVs show strong Fe Kα lines in their X-ray spectra (e.g., Norton et al. 1991; De Martino et al. 2008; Byckling et al. 2010; Kuulkers et al. 2010; Mukai 2017). We do not find any evidence for such a line in the X-ray spectra of CX1 (even after merging the spectra). However, it is possible that this is due to poor statistics in the X-ray spectra or full ionization of Fe in the emitting region.

The X-ray spectrum of CX1 is similar to many LMXBs containing black holes or neutron stars. In the X-ray binary scenario for CX1, the bright X-ray detections in 2003 and 2016 could be explained as accretion outbursts. However, most LMXBs spend most of their outburst at X-ray luminosities above $10^{35}$ erg s$^{-1}$. This would mean that both 2003 and 2016 observations were taken during the early rise or late decay of potential outbursts from CX1, which seems unlikely.

One class of X-ray transients known as very faint X-ray transients (VFXTs) undergo faint outbursts, with X-ray luminosity peaking between $\sim 10^{34}$ and $10^{36}$ erg s$^{-1}$ during the outburst. Many such transients have been found in surveys of the Galactic center (Muno et al. 2005; Sakano et al. 2005; Wijnands et al. 2006; Degenaar & Wijnands 2009a, 2009b), and a few are found in globular clusters (e.g., Heinke et al. 2009; Arnason et al. 2015). While most of these systems show short outbursts (timescale of hours to days), there are some that show quasi Persistent outbursts (e.g., Armas Padilla et al. 2013). The nature of these systems and their low accretion rates are not fully understood, but they are likely to be a somewhat heterogeneous population. It has been suggested that VFXTs might be “period-gap” systems (Maccarone & Patruno 2013) or that some of them might be unidentified tMSPs (Heinke et al. 2014), and recently Hofmann et al. (2018) identified a new transient intermediate polar system as a VFXT. The behavior of CX1 in X-rays is similar to some VFXTs.

It is possible that black hole LMXBs in globular clusters show different outburst behavior than field systems due to their dynamical formation, including systems that show only faint short outbursts (Knevitt et al. 2014). This is certainly a fascinating possibility for CX1. However, confirmed accreting neutron stars substantially outnumber even candidate accreting black holes in globular clusters (e.g., Bahramian et al. 2014, table 510), and the large MSP population of Ter 5 means that it has hosted many neutron star LMXBs, so in the absence of other constraints, a neutron star accretor should be considered more likely.

### 5. Conclusions

We conclude that all available evidence is consistent with the hypothesis that Terzan 5 CX1 is a tMSP that was in a low-level accretion state in 2003, transitioning to a pulsar state some time before 2009. The high density of observations between 2009 and 2014 suggests it was largely or entirely in the pulsar state before 2009. The high density of observations between 2009 and 2014 suggests it was largely or entirely in the pulsar state during this time. By 2016, it had returned to a low-level accretion state (Figure 2).

Extending the baseline further, early Chandra observations, obtained during the outburst of the low-mass X-ray binary EXO 1745-248 (Heinke et al. 2003a), provide a 3σ upper limit of $L_X < 4.7 \times 10^{32}$ erg s$^{-1}$. This suggests that in 2000 the system may have been in the fainter pulsar state, indicating the possibility of an additional transition between 2000 and 2003.

Given the extensive radio pulsar timing data set available for Terzan 5, the best way to confirm CX1 as a tMSP would be to detect the source as a radio pulsar using data obtained during a putative pulsar phase (e.g., searching the data obtained between 2009 and 2014). Such searches are ongoing, and new pulsars

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10 https://bersavosh.github.io/research/gc_lmxbs.html
are being detected in older Terzan 5 timing observations using refined search techniques (e.g., Cadelano et al. 2018). It also suggests that deep, well-separated epochs of X-ray data should be searched for other globular clusters, and could reveal the presence of additional tMSPs.

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**Facilities:** *Chandra*(ACIS), VLA, HST (ACS).

**Software:** AIPS (Greisen 2003), Astropy (Astropy Collaboration et al. 2013), CIAO (Fruscione et al. 2006), HEASOFT (Nasa High Energy Astrophysics Science Archive Research Center 2014), Matplotlib (Hunter 2007), PyRAF (Science Software Branch at STScI 2012), SAOImage DS9 (Smithsonian Astrophysical Observatory 2000), Xspec (Arnaud 1996).

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