Computational reconstruction of pacemaking and intrinsic electroresponsiveness in cerebellar Golgi cells

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The Golgi cells have been recently shown to beat regularly in vitro (Forti et al., 2006, J. Physiol., 574, 711–729). Four main currents were shown to be involved, namely a persistent sodium current (\(I_{Na-p}\)), an h current (\(I_h\)), an SK-type calcium-dependent potassium current (\(I_{K_{AHP}}\)), and a slow M-like potassium current (\(I_{K_{slow}}\)). These ionic currents could take part, together with others, also to different aspects of neuronal excitability like responses to depolarizing and hyperpolarizing current injection. However, the ionic mechanisms and their interactions remained largely hypothetical. In this work, we have investigated the mechanisms of Golgi cell excitability by developing a computational model. The model predicts that pacemaking is sustained by subthreshold oscillations tightly coupled to spikes. \(I_{Na-p}\) and \(I_{K_{slow}}\) emerged as the critical determinants of oscillations. \(I_h\) also played a role by setting the oscillatory mechanism into the appropriate membrane potential range. \(I_{K_{AHP}}\), though taking part to the oscillation, appeared primarily involved in regulating the ISI following spikes. The combination with other currents, in particular a resurgent sodium current (\(I_{Na-r}\)) and an A-current (\(I_{K_{A}}\)), allowed a precise regulation of response frequency and delay. These results provide a coherent reconstruction of the ionic mechanisms determining Golgi cell intrinsic electroresponsiveness and suggests important implications for cerebellar signal processing, which will be fully developed in a companion paper (Solinas et al., 2008, Front. Neurosci. 2:4).

Keywords: Golgi cell, pacemaking, cerebellum, granular layer, modeling

INTRODUCTION

Oscillations are one of the most prominent features of brain activity (Buzsaki, 2006). Numerous neurons in the brain show pacemaker activity exploiting a variety of ionic mechanisms (Hutcheon and Yarom, 2000; Koch, 1999; Llinás, 1988; Yamada et al., 1998). In a well-known case represented by relay thalamic neurons (RTN), large low-voltage oscillations are generated by the interplay of h-current with low-threshold activated (LVA) calcium channels (Huguenard and McCormick, 1992; McCormick and Huguenard, 1992; McCormick and Pape, 1990). The h-current determines the depolarizing ramp leading to LVA channel activation and to the calcium spike, which triggers a spike bursts before inactivating and concluding the cycle. These pacemakers cannot be easily perturbed and drive fast sodium spike bursts. Some other, probably more diffused pacemaker neurons, show smaller oscillations in the subthreshold membrane potential region generating a regular low-frequency beating, usually in the theta-frequency band. These include several neuronal types in the basal ganglia and brain cortex (Bennett et al., 2000; Dickson et al., 2000a; Surmeier et al., 2005; Wilson and Callaway, 2000). The mechanisms are thought to employ primarily persistent sodium currents in combination with some slow repolarizing mechanism, although oscillations driven by high-voltage activated calcium channels have also been reported. The variety of pacemaker mechanism may be related to the additional behaviors the neuron is able to express. It is therefore important to determine how a specific neuron generates pacemaking and how this combines with other aspects of electroresponsiveness.

The Golgi cell (Golgi, 1883) is an inhibitory interneuron regulating information flow in the granular layer circuit of cerebellum (Eccles et al., 1967; Marr, 1969), which has been recently shown to beat regularly in vitro (Forti et al., 2006, see also Dieudonné, 1998). Electrophysiological and pharmacological analysis (Forti et al., 2006) showed that four main currents were involved, namely a persistent sodium current (\(I_{Na-p}\)), an h current (\(I_h\)), an SK-type calcium-dependent potassium current (\(I_{K_{AHP}}\)), and a slow M-like potassium current (\(I_{K_{slow}}\)). However, the precise role of these currents was unclear. In particular, both \(I_{Na-p}\) and \(I_h\) may be able to drive the oscillation and the repolarizing action of \(I_{K_{AHP}}\) and \(I_{K_{slow}}\) may also be redundant.

As a first step for investigating the mechanisms of Golgi cell excitability, we have developed a computational model incorporating state-of-the-art knowledge on Golgi cell ionic channels and electro-
responsiveness. The model predicted that pacemaking was sustained by subthreshold oscillations generated by $I_{h}$ and $k_{s}$ was less important but maintained the model into the appropriate membrane potential range for oscillation. The oscillations were tightly coupled to spikes thereby causing a prominent activation of $k_{s}$ and a stabilization of the pacemaker cycle. The combination with other currents, in particular a resurgent sodium current ($I_{h}$) and an A-current ($k_{a}$), allowed a precise regulation of response frequency and delay. These results provide a coherent reconstruction of the ionic mechanisms determining Golgi cell intrinsic electroresponsiveness. The implications of this complex ionic complement will be further clarified in a companion paper (Solinas et al., 2007), in which $k_{s}$ and $k_{l}$ will assume the additional and more specific roles of determining phase-reset and theta-frequency resonance.

### MATERIALS AND METHODS

This work presents a combined experimental and modeling analysis of Golgi cell electroresponsiveness. In order to develop a comprehensive hypothesis on how the Golgi cell generates its excitatory response, an extensive re-analysis of patch-clamp recordings and pharmacological data obtained in a previous paper (Forti et al., 2006) has been complemented with a computational model. The model allowed us to investigate neuronal behaviors hard to resolve experimentally, like the apparent close coupling of spikes with intrinsic oscillatory mechanisms and the effect of persistent Na$^+$ currents. Moreover, while conductance measurements and dynamic-clamp have been successfully applied to electrotonically compact neurons (e.g., in isolated Purkinje cells: Raman and Bean, 1997), the large clamp escape in Golgi cells in slices (data not shown) prevents the application of these methods.

### Golgi cell modeling

A computational model of rat Golgi cell electroresponsiveness was constructed using the NEURON simulator (version 5.9; Hines and Carnevale, 2001). The model consisted of five compartments allowing for a minimal description of the Golgi cell electrotonic structure. The soma was connected to an “electrode” to reproduce realistic current-clamp conditions. All voltage- and Ca$^{2+}$-dependent mechanisms (see below) were placed in the somatic compartment. With this approach, the model reproduced satisfactorily basic aspects of Golgi cell electroresponsiveness elicited by somatic current injection. It should also be noted that, although differences have been reported at the histochemical level (Geurts et al., 2001; Simat et al., 2007), the basic electroresponsive properties were homogeneous in a large majority of Golgi cells (Forti et al., 2006, see also figures in this paper). Thus, we have reconstructed a “canonical” Golgi cell model which simulates the electrophysiological behavior for this major Golgi cell subgroup and lays within the scatter of physiological parameter values. The model was calibrated on the cellular response corrected for the liquid-junction potential.

The model included 12 voltage-dependent conductances placed into the soma (Table 1). Gating kinetics were corrected using a $Q_{10}$ for $I_{h}$ to $V_{m}$ activation (Destrée et al., 1994), and a $Q_{10} = 3$ for all the other currents according to the relation $Q_{10}^{(T_{m} - T_{ref})/T_{ref}}$ (Gutfrndet et al., 1995) to account for temperature differences between experimental recordings and the model (23ºC). Nernst equilibrium potentials were calculated from ionic concentrations used in current-clamp recordings. Only the Ca$^{2+}$ Nernst potential was updated during simulation. The maximum density voltage-dependent conductances were regulated to fit the Golgi cell responses to various stimulations (Achard and De Schutter, 2006; Traub and Llinás, 1979; Traub et al., 1991; Vanier and Bower, 1999).

Ionic currents were modeled following the Hodgkin and Huxley formulation of ion channel gating using mathematical methods reported previously (D’Angelo et al., 2001; Nieus et al., 2006). Voltage was obtained as the time integral of the equation (Yamada et al., 1998)

$$dV/dt = -I_{m}/C_{m} = \left[\sum g_{i} (V - V_{i}) + i_{inj}\right]$$

where $I_{m}$ is membrane potential, $C_{m}$ membrane capacitance, $g_{i}$ are ionic conductances, and $V_{i}$ reversal potentials (the subscript $i$ indicates different channels), and $i_{inj}$ is the injected current. Hodgkin and Huxley is simulated using a Markov gating scheme reproducing the dynamics reported by Hirschberg et al., (1998), while all other membrane conductances were represented using Hodgkin–Huxley-like models (Hodgkin and Huxley, 1952) of the type

$$g_{i} = G_{max} \times (V_{th} - V_{i})$$

where $G_{max}$ is the maximum ionic conductance, $V_{th}$, and $V_{i}$ are state variables (probabilities ranging from 0 to 1) for a gating particle, and $V_{th}$ is the such of gating particles in ionic channel $i$. $x$ and $y$ (the suffix $i$ omitted) were related to the first order rate constants $\alpha$ and $\beta$ by the equations

$$x_{\infty} = \frac{\alpha_{x}}{\beta_{x}}$$

$$y_{\infty} = \frac{\alpha_{y}}{\beta_{y}}$$

$$\tau_{x} = \frac{1}{\alpha_{x} + \beta_{x}}$$

$$\tau_{y} = \frac{1}{\alpha_{y} + \beta_{y}}$$

where $\alpha$ and $\beta$ are functions of voltage. The equations used to parameterize $\alpha$ and $\beta$ and the state variables $x_{\infty}$, $\tau_{x}$, $y_{\infty}$, $\tau_{y}$ for different ionic channels are shown in Table 1. The state variable kinetics were

$$\frac{dx}{dt} = (x_{\infty} - x)/\tau_{x}$$

$$\frac{dy}{dt} = (y_{\infty} - y)/\tau_{y}$$

The intracellular Ca$^{2+}$ concentration, $[Ca^{2+}]_{i}$, was calculated through the equation

$$\frac{d[Ca^{2+}]}{dt} = -\frac{I_{Ca}}{(2F \times Ad)} - \frac{1}{\beta_{Ca}} ([Ca^{2+}]_{i} - [Ca^{2+}]_{eq})$$

where $d$ is the depth of a shell adjacent to the cell surface of area $A$, $\beta_{Ca}$ determines the loss of Ca$^{2+}$ ions from the shell approximating the effect of fluxes, ionic pumps, diffusion, and buffers (De Schutter and Smolen, 1998; McCormick and Huguenard, 1992; Traub and Llinás, 1979), and $[Ca^{2+}]_{eq}$ is the resting [Ca$^{2+}$].

### Model mechanisms

**Passive properties and model structure.** In cerebellar Golgi cells, charge displacement during voltage steps followed a tri-phasic trajectory (see Figure 1). According to Dieudonné (1998), the fastest transient was related to the soma, the intermediate to the dendrites, and the slowest to the axon. According to proportions reported by Dieudonné (1998) the capacitances of soma, dendrites, and axon were set to 23, 32, and 90 pF, respectively, summing up to a total cell capacitance of 145 pF (Forti et al., 2006). The specific axial resistance of axon and dendrites was set to 100 $\Omega \cdot cm$ (Roth and Häusser, 2001), the specific membrane resistance was set to 47.6 k$\Omega \cdot cm^{-2}$ (corresponding to the experimental determination of leakage reported below) and specific membrane capacitance was set at 1 $\mu$F/cm$^{2}$. Compartment equivalent to the complex structure of dendrites and axons (Koch, 1999; Yamada et al., 1998) were generated to comprise a spherical soma (diameter = 27 $\mu$m), a single axonal compartment (1200 $\mu$m long and 2.4 $\mu$m thick), and three dendritic compartments (each 113 $\mu$m long and 2.4 $\mu$m thick). With these parameters, input capacitance and input resistance in the model equaled those of Golgi cells and the model could appropriately reproduce current transients recorded in vivo in the subthreshold region (Figure 1). The micropipette was simulated as a passive cable with sealed end, null capacitance, infinite
Table 1. Voltage-dependent conductance parameters.

| Conductance state variables | $n$ | $g_{\text{max}}$ (S/cm$^2$) | $V_{\text{m}}$ (mV) | $\alpha$ (s$^{-1}$) | $\beta$ (s$^{-1}$) | Reference |
|-----------------------------|-----|--------------------------|---------------------|-----------------|-----------------|----------|
| $E_{\text{Na}}$-t | Activation | 3 | 0.048 | 77.39 | 0.417(V + 35)/[1−exp(−(V + 35)/10)] | 16.68 exp(−0.055(V + 60)) | Modified from Nieus et al., 2006 |
| | Inactivation | 1 | | 0.292 exp(−0.3(V + 60)) | | |
| $E_{\text{Na}}$-r | Activation | 1 | 0.0017 | 77.39 | [1.11−68.5(V + 6.481)/exp(−(V + 6.481)/6.8 − 1)] × 10$^{-4}$ | [66−21.7(V + 54)/exp((V + 54)/0.11) − 1] × 10$^{-3}$ | |
| | Inactivation | 1 | | 0.443 × exp(−(V + 90)/62.5) | | |
| $E_{\text{Na}}$-p | Activation | 1 | 1.9 × 10$^{-4}$ | 77.39 | 0.421(V + 50)/[1−exp(−(V + 50)/5)] | −0.287(V + 50)/[1−exp(−(V + 50)/5)] | |
| | Inactivation | 1 | | | | |
| $E_{\text{K}}$ | Activation | 4 | 0.032 | −94.6 | 0.062(V + 36)/[1−exp(−(V + 36)/10)] | 0.78 exp(−(V + 46)/80) | Hirschberg et al., 1998 |
| | Inactivation | 1 | | | | |
| $E_{\text{Ca}}$ | Activation | 3 | 0.008 | −94.6 | 0.62/[1+exp(−(V + 19.17)/23.32)] | 0.126 exp(−(V + 28.28)/19.47) | Destenche et al., 1994 |
| | Inactivation | 1 | | 0.026/[1+exp(−(V + 121)/12.84)] | | |
| $E_{\text{K}}$ | Activation | 1 | 0.009 | −94.6 | 3.2/[1+0.005 exp(−(V + 10)/11.7)] | 0.46/[1+exp(−(V + 88.8)/8.4)] | |
| | Inactivation | 1 | | | | |
| $E_{\text{K-slow}}$ | Activation | 1 | 0.001 | −94.6 | 0.0037 exp((V + 40)/40) | 0.0037 exp(−(V + 40)/20) | |
| | Inactivation | 1 | | | | |
| $E_{\text{Ca-HVA}}$ | Activation | 2 | 4.6 × 10$^{-4}$ | 136.3$^*$ | 0.0687 exp(0.063(V + 39)) | 0.115 exp(−0.039(V + 28.66)) | Santoro et al., 2000 |
| | Inactivation | 1 | | 1.8 × 10$^{-3}$ exp(−(V + 58)/18) | 1.8 × 10$^{-3}$ exp(−(V + 58)/83) | | |
| $E_{\text{Ca-LVA}}$ | Activation | 2 | 2.5 × 10$^{-4}$ | 136.3$^*$ | $x_c = 1/(1+exp(−(V + 82)/7.4))$ | $\tau_c = (3+1/[exp(−(V + 37)/10)] + \exp(−(V + 112)/15))/0.85$ | Reconstructed from Santoro et al., 2000 |
| | Inactivation | 1 | | $y_c = 1/(1+exp(−(V + 90)/5))$ | $\tau_y = (85+1/[exp(−(V + 58)/4)] + \exp(−(V + 417)/50))/0.9$ | | |
| $E_{\text{K-AHP}}$ | 0.038 | −94.6 | $c_1 \leftrightarrow c_2 \alpha = 200$ ms$^{-1}$ m$^{-1}$; $\beta = 0.08$ ms$^{-1}$ | $c_2 \leftrightarrow c_4 \alpha = 160$ ms$^{-1}$ m$^{-1}$; $\beta = 0.08$ ms$^{-1}$ | $c_3 \leftrightarrow c_4 \alpha = 80$ ms$^{-1}$ m$^{-1}$; $\beta = 0.2$ ms$^{-1}$ | $c_3 \leftrightarrow c_5 \alpha = 0.16$ ms$^{-1}$; $\beta = 1$ ms$^{-1}$ | Hirschberg et al., 1998 |
| | | | $c_1 \leftrightarrow 3c_2+c_1 \leftrightarrow c_4$ | | | $c_4 \leftrightarrow 2c_5 \alpha = 1.2$ ms$^{-1}$; $\beta = 0.1$ ms$^{-1}$ | |
| $E_{\text{HCN}1}$ | 1 | 5 × 10$^{-5}$ | −30 | \begin{align*} r &= 0.002(1+10^{-6}) + 0.97 \\
 h_{\text{max}} &= r/[1+exp(−(V + 82.5)/0.11))] \\
 h_{\text{down}} &= [1−r]/[1+exp(−(V + 82.5)/0.11))] \\
 \tau_{\text{down}} &= \exp(0.014V−14.07)/2.3 \\
 \tau_{\text{down}} &= \exp(0.014V−14.07)/2.3 \\
 \tau_{\text{down}} &= \exp(0.014V−14.07)/2.3 \\
 \end{align*} | \begin{align*} r &= 0.002(1+10^{-6}) + 0.97 \\
 h_{\text{max}} &= r/[1+exp(−(V + 91.9)/0.16))] \\
 h_{\text{down}} &= [1−r]/[1+exp(−(V + 91.9)/0.16))] \\
 \tau_{\text{down}} &= \exp(0.015V+15.3)/2.3 \\
 \tau_{\text{down}} &= \exp(0.015V+15.3)/2.3 \\
 \end{align*} | \begin{align*} r &= 0.002(1+10^{-6}) + 0.97 \\
 h_{\text{max}} &= r/[1+exp(−(V + 91.9)/0.16))] \\
 h_{\text{down}} &= [1−r]/[1+exp(−(V + 91.9)/0.16))] \\
 \tau_{\text{down}} &= \exp(0.015V+15.3)/2.3 \\
 \tau_{\text{down}} &= \exp(0.015V+15.3)/2.3 \\
 \end{align*} | Reconstructed from Santoro et al., 2000 |

The table reports the equations used to calculate $\alpha$, $\beta$, $\alpha$, and $\beta$ (see Equations (3) and (4)) for the membrane conductances used in the model (in a few cases $x_{\infty}$ and $y_{\infty}$ are reported). Only the SK channel was built using Markov multistate transition model in order to use the data available in literature (Santoro et al., 2000). The number of gating particles ($n$), maximum conductance ($g_{\text{max}}$), and ionic reversal potential ($V_r$ resting value for Ca$^{2+}$) used to calculate ionic currents are also shown. The opening and closing rates are corrected for temperature (23°C).

**Voltage-dependent mechanisms.** These were based on experimental observations including current-clamp (Forti et al., 2006) and preliminary voltage-clamp recordings. A first current set was needed to generate action potentials. This consisted in a standard description of $h_{\text{Na}}$, and $h_{\text{K}}$, and was completed through the introduction of $I_{\text{Ca-HVA}}$ and $h_{\text{Ca}}$, which improved the fast phase (1−2 ms following the upstroke) of spike afterhyperpolarization (AHP) preventing Na$^+$ channel inactivation and discharge blockade at high frequency. $h_{\text{K}}$ regulated the first spike delay and $h_{\text{Na}}$, intensified high-frequency discharge. $I_{\text{Ca-LVA}}$ was needed to enhance rebound depolarization. A second current set was needed to generate pacemaking and, according to previous experimental analysis (Forti et al., 2006), included $I_{\text{K-AHP}}$, $h_{\text{K-slow}}$, and $h_{\text{K-AP}}$. The procedures include a range of wall resistance and an axial resistance of 25 MΩ equivalent to “access resistance” (Neher and Sakmann, 1995).

**Leakage.** The leakage current ($I_{\text{leak}}$) was measured in the subthreshold region from VC recordings after application of the following ionic channel blockers (in mM): TEACl (40), BaCl$_2$ (2), CdCl$_2$ (0.2), CoCl$_2$ (2), TTX (0.001), Apamin (0.05), Linopirdine (0.04), D72728S (0.01−0.1). All drugs were applied through general bath perfusion as described in Forti et al., 2006. The currents obtained from voltage-step injection yielded a linear $I$-$V$ plot (range −120 to −60 mV) maximum conductance of 21 μS/cm$^2$ and reversal potential at −60 mV (data not shown).
adopted to reproduce ionic channel gating (Hodgkin and Huxley, 1952; Neher and Sakmann, 1995) and set current density (Achard and De Schutter, 2006; Traub and Llinás, 1979; Traub et al., 1991; Vanier and Bower, 1999) are reported below.

**I_p**. The presence of \( I_p \) in Golgi cells was reported by Dieuédonné (1998) and confirmed by Forti et al. (2006). \( I_p \) in Golgi cells is manifest during the application of hyperpolarizing current steps causing sagging inward rectification with both a rapid and a slow phase. Slow and fast response components were reproduced by two voltage-dependent mechanisms corresponding to HCN1 and HCN2 currents, respectively (Santoro et al., 2000; see also Chan et al., 2004; Dickson et al., 2000b; Luthi and McCormick, 1998; McCormick and Pape, 1990; McCormick and Huguenard, 1992). After regulating current densities to match the sag and the steady-state level of hyperpolarizing responses, the model generated a robust rebound depolarization, as observed experimentally (a reinforcement of the early rebound phase was obtained with \( I_{Ca-LVA} \)), see below.

**I_{Na-p}**. \( I_{Na-p} \) was matched to the \( I-V \) curve reported by Forti et al. (2006). \( I_{Na-p} \) kinetics were derived from the model previously developed for cerebellar granule cells (D’Angelo et al., 2001; Magistrelli et al., 2006; Nieus et al., 2006) and pyramidal cells (Gutfreund et al., 1995).

**I_{K-stow}**. The presence of a slow repolarizing retigabine-sensitive M-like current \( I_{K-stow} \) in Golgi cells was suggested by pharmacological observations (Forti et al., 2006). The \( I_{K-stow} \) model was obtained from previous models (D’Angelo et al., 2001; Gutfreund et al., 1995) and maximum conductance was set in proportion to \( I_{Na-p} \) (Nieus et al., 2006). \( I_{K-stow} \) contributed to shape the ISI and spike frequency adaptation.

**I_{Ca-AHP}**. The presence of a slow \( Ca^{2+} \)-dependent repolarizing apamin-sensitive \( K^- \) current (encoded by SK-type channels; for a review see Bond et al., 2005; Stocker, 2004) sustaining the slower phase (peaking 30–40 ms after spike peak) of the spike AHP in Golgi cells was indicated by pharmacological observations (Forti et al., 2006). Although, in rat Golgi cells, mRNAs encoding the SK3 subtype are particularly abundant (Stocker and Pedarzani, 2000), the corresponding current in the model, \( I_{Ca-AHP} \), was simulated using the available SK2 model reported by Hirschberg et al. (1998). \( I_{Ca-AHP} \) was activated by \( Ca^{2+} \) entering through HVA \( Ca^{2+} \) channels (Marrion and Tavalin, 1998; Stocker, 2004). After regulating intracellular \( Ca^{2+} \) dynamics, \( I_{Ca-AHP} \) was tuned in order to fit the slow AHP phase and spike frequency adaptation.

**I_{Ca-LVA}**. Preliminary observations suggest that Golgi cells express LVA calcium channels reinforcing rebound excitation on return from hyperpolarization (Forti, Cesana and D’Angelo, unpublished observations). Indeed, the model could not properly reproduce the initial reinforcement of rebound firing with \( I_p \) alone. The \( I_{Ca-LVA} \) model was taken from Destexhe et al., (1994) and current density adjusted to appropriately reinforce rebound excitation. Due to its voltage-dependent inactivation, \( I_{Ca-LVA} \) did not contribute to any other excitatory process of the Golgi cell.

**Spike-related currents.** The complement of spike-related currents was reconstructed with mechanisms taken from a previous granule cell model (D’Angelo et al., 2001; Nieus et al., 2006). The density of the transient \( Na^- \) current (\( I_{Na} \)) was adjusted to reproduce the action potential threshold and overshoot. A preliminary investigation revealed that, as well as granule cells, Golgi cells also express a proportionate resurgent \( Na^- \) current (\( I_{Na-r} \)) component (Magistrelli, Castelli and D’Angelo, unpublished observations). Its insertion into the model was required to enhance the initial doublet in response to depolarization raising the maximum firing rate into the physiological range measured in vitro (up to ~200 Hz, Forti et al., 2006). The delayed rectifier \( K^- \) current, \( I_{K-V} \), was needed to generate spike repolarization and balanced with \( I_{K-C} \) to match the spike shape. The requirement of a \( Ca^{2+} \)-dependent voltage-dependent \( K^- \) current, \( I_{K-C} \), emerged from the inability of \( I_{K-V} \) to fully account for the fast phase of spike AHP and to guarantee \( Na^- \) channel re-priming at high discharge frequency. \( I_{K-A} \) was needed to regulate spike delay in response to depolarizing current injection and during rebound excitation.

**Ca^{2+} dynamics.** \( Ca^{2+} \) dynamics were adapted to yield \( Ca^{2+} \) transients of about 1 \( \mu \)M following spike-dependent opening of HVA channels (Gabbiani et al., 1994). Parameters used in Equation (6) \((d = 200 \text{ nm}, \beta_{Ca} = 1.3, \text{ and } [Ca^{2+}]_o = 50 \text{ nM})\) were similar to those used in other neurons (e.g., De Schutter and Smolen, 1998; D’Angelo et al., 2001; Traub and Llinás, 1979; Traub, 1982). With this asset, there was calcium accumulation during high-frequency discharge but not when the model fired at its intrinsic pacemaker rhythm. The computed intracellular \( Ca^{2+} \) concentration \([Ca^{2+}]_i\), activated \( I_{K-C} \) just after the spike upstroke and generated the fast AHP. \( I_{Ca-AHP} \) was activated using 1/3 of the computed \([Ca^{2+}]_i\), value on the assumption that local concentration is smaller on SK than on channels supporting \( I_{K-C} \) (see Stocker (2004) for a discussion on effective local \([Ca^{2+}]_i\) concentrations). This seems reasonable in view of the higher \([Ca^{2+}]_i\), affinity of SK channel subtypes with respect to channels supporting the fast spike AHP (reviewed in Stocker, 2004). The slow dynamics of \( I_{Ca-AHP} \), which contributed to the slower phase of the spike AHP, was accounted for by slow channel closing kinetics rather than by the permanence of elevated \( Ca^{2+} \) levels. \( I_{Ca-AHP} \) activation was enhanced during high-frequency repetitive discharge due to the associated \( Ca^{2+} \) build-up, thereby determining spike-frequency adaptation (see Figure 4). It should be noted that, in the model, \( I_{Ca-LVA} \) was not allowed to contribute to the \( Ca^{2+} \) pool controlling \( I_{K-C} \) and \( I_{Ca-AHP} \), since such a coupling would unrealistically depress rebound excitation (data not shown).
Membrane noise in vitro. In order to reproduce the ISI irregularity observed during Golgi cell recordings in vitro (CV_{ISI} = 0.13, median value from loose cell-attached recordings, Forti et al., 2006; CV_{ISI} = 0.16, median value from WCR, this study), Gaussian noise was added to the model in selected simulations aimed at reproducing the effect of apamin application (cf. Figures 2E and 5C). The NEURON random number generator was used to produce a sequence of numbers cast from a Gaussian distribution with 0 mean and unitary variance generating a sequence lacking temporal correlation. The random numbers were scaled by a 32 pA coefficient yielding a noise current injected into the soma. In the canonical model, this produced a CV_{ISI} = 0.16.

The Golgi cell model NEURON code is available from ModelDB (senselab.med.yale.edu/ModelDB/).

Data analysis

The processing of experimental and simulated traces was automated using dedicated scripts written in MATLAB (MathWorks, Natick, MA, USA). As in Forti et al., (2006), spike threshold was detected at the point where the depolarization rate reached 5 mV/ms. Spike maximum and minimum (AHP trough) were computed and the AHP rise-time was estimated as the time needed to reach the AHP trough from threshold crossing during the down phase of the spike. The sag amplitude was calculated as the difference between the minimum and steady-state value of membrane potential during responses to hyperpolarizing currents. Firing rate adaptation in response to depolarizing currents was calculated as the ratio of the steady-state and initial discharge frequency (i.e., the inverse of the last and first ISI).

The match of model to experimental data was evaluated by testing the null hypothesis that their difference was a random sample from a normal distribution (Lilliefors test) with mean equal to 0 and unknown variance (Student’s t-test). All data are reported as mean ± SD.

RESULTS

The Golgi cell was recently proposed to express a characteristic set of ionic currents (comprising Na-p, h, K-AHP, and K-slow; Forti et al., 2006). Each one of these is endowed with unique gating and kinetics properties (see for a comprehensive description Koch, 1999; Yamada et al., 1998): Na-p activates rapidly with depolarization in the subthreshold region, h slowly deactivates with depolarization, K-AHP activates briskly upon Ca^2+ entry and remains activated for tens of milliseconds, and K-slow activates slowly and progressively in the subthreshold region during sustained depolarization.

In order to elaborate a hypothesis on how these ionic mechanisms interact, the electrophysiological and pharmacological analysis of intrinsic excitability was assisted by a computational model incorporating prototypical ionic current properties (see Methods, Figure 1 and Table 1).

Modeling the basic properties of Golgi cell electroresponsiveness

The basic properties of Golgi cell electroresponsiveness in acute brain slices (henceforth called in vitro) were inferred from the previous paper of Forti et al., (2006) and from novel WCR data, and could be summarized as follows. (i) Golgi cells paced rhythmically at a frequency between 1 and 8 Hz at room temperature (Figure 2A). In the pacemaking regime, spikes were followed by a protracted AHP (Figure 2B). The average ISI distribution obtained from different cells in WCR showed a prominent peak around 200 ms (Figure 2C). There was no clear correlation between the

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Since the inspection of data plots in Figures 1–4 did not allow to discern subclasses of Golgi cells on the basis of pacemaking frequency, ISI shape, rebound excitation or discharge frequency, a “canonical” model was generated to simulate the average electrophysiological behavior. The model was able to quantitatively reproduce all the properties summarized above, as demonstrated in Figures 1–4. The average pacemaker frequency coincided with the mode of the experimental data distribution, and model predictions about spike shape, inward rectification, rebound excitation, firing delay, and frequency adaptation during depolarization fell within the experimental data distributions with $p > 0.12$ (t-test). It should be noted that the AHP at the end of a depolarization-induced discharge deviated from the shape observed in neurons (Figure 4A) suggesting that the model does not completely represent the $\text{Ca}^{2+}$ dynamics developing over longer time-scales. This is, however, a minor problem in the present context, since the physiological implications of AHP are considered in response to short depolarization and the model could indeed precisely predict the subsequent silent phase (see Solinas et al., 2007).

Testing the model versus pharmacological challenges

The adequacy of the model in terms of intrinsic mechanisms was further tested by comparison with pharmacological observations using specific ionic channel blockers (cf. Forti et al., 2006). A 50% $I_{\text{h}}$ reduction, simulating sub-maximal block by ZD7288, reduced the pacemaker frequency (73% of control) and slowed down the ISI trajectory (Figure 5A), similar to in vitro observations. As expected from their gating properties reported in vitro, 50% $I_{\text{h}}_{\text{a}}$ reduction arrested pacemaking (Figure 5A). The action of retigabine, an agonist of M-type currents which reduced firing frequency and altered the ISI during current injection in vitro, was matched by a $-6.5 \text{ mV}$ shift in $I_{\text{h}}_{\text{a}}$ activation (Tatulian et al., 2001) (Figure 5B). Finally, since the most evident effect of the SK channel blocker apamin was to reduce pacemaker regularity, the role of $I_{\text{AHP}}$ in the model was investigated after adding noise, to reproduce the standard deviation of membrane potential during the ISI (Figure 5C; see Methods for details). The simulated pacemaking thus acquired a slight irregularity ($\text{CV} = 0.16$), which was markedly increased by $I_{\text{AHP}}$ blockade ($\text{CV} = 0.36$), as observed in the experiments in vitro. It is thus probable that the AHP, by temporarily displacing membrane potential below threshold, prevents the perturbing action of stochastic channel gating and spontaneous synaptic currents in the subthreshold regime. In summary, simulations show that the model predictions for the effects of $I_{\text{h}}_{\text{a}}, I_{\text{h}}_{\text{slow}}, I_{\text{AHP}}$, and $I_{\text{h}}$ on pacemaking agree with pharmacological analysis in vitro.

Model currents and intracellular calcium

The model currents and intracellular calcium changes during Golgi cell activity are shown in Figure 6. The currents dominating the ISI during pacemaking were small (<5%) compared to those generating the spikes. During the ISI, $I_{\text{h}}_{\text{a}}$ showed a progressive increase, while $I_{\text{h}}$ remained almost constant. $I_{\text{h}}_{\text{AHP}}$ was driven by the spike, arose rapidly and then decreased toward the middle of the ISI. $I_{\text{h}}_{\text{slow}}$ was driven by depolarization in the second part of the ISI and showed a progressive increase. It should be noted that $I_{\text{Ca-HVA}}$ activated during the spike upstroke. During pacemaking, intracellular $\text{Ca}^{2+}$ returned to basal levels within an ISI, but during high-frequency bursts $\text{Ca}^{2+}$ could accumulate and thereby intensify and protract $I_{\text{h}}_{\text{AHP}}$ activation. Spike-frequency adaptation during discharge elicited by current injection was caused by $I_{\text{h}}_{\text{AHP}}$ and $I_{\text{h}}_{\text{slow}}$. Rebound excitation was determined by $I_{\text{h}}$ deactivation and intensified by $I_{\text{Ca-LVA}}$. Certain currents accurately regulated the responses to current injection (Figure 7). During depolarization, $I_{\text{A}}$ delayed the first spike and $I_{\text{h}}_{\text{a}}$ enhanced the initial doublet. As well as during depolarization, $I_{\text{A}}$ contributed to regulate the delay of rebound excitation (data not shown).
Robustness and the canonical neuron

The simulations we have presented till now lead to two conclusions. (i) Although the currents introduced into the model (on the basis of experimental observations) controlled characteristic aspects of electroresponsiveness (see also Solinas et al., 2007), they could have multiple and redundant effects. (ii) The canonical Golgi cell model is one that succeeds in approximating the behavior of the experimental population of Golgi cells. The robustness of this model should be tested by checking its ability to reproduce the variability of behaviors measured in different cells when varying model parameters. We therefore measured robustness of the model by assessing its ability to maintain typical Golgi cell characteristics when key parameters were varied. This aspect is shown in Figure 8, in which certain current densities or the time constant of $[\text{Ca}^{2+}]_i$ are modified in turn. The analysis in Figure 8 confirms the relative importance of different currents in generating Golgi cell behavior reported in Figure 5. The model was quite sensitive to changes in $I_{\text{Na-p}}$ conductance, followed by changes in $I_{\text{K-slow}}$. Calcium dynamics, determined by the conductance of $I_{\text{Ca-HVA}}$ and the decay constant $\beta_{\text{Ca}}$ were also relatively important, while the model was quite insensitive to changes in $I_h$ conductance. An example was presented in Figure 3, in which some parameters were modified to adapt the canonical model to mimic an individual Golgi cell response in detail. Notice that a more extended robustness assessment (Achard and De Schutter, 2006; Druckmann et al., 2007) would imply also testing the effect of conjoint variations of free parameters. This is a complex task that may be addressed in future research.

Coupling of subthreshold membrane potential oscillations with spikes

Given the good response of the model over a wide range of experimental behaviors, it could be used to predict the mechanism of pacemaking. In several neurons (e.g., in amygdaloid neurons (Papae and Driesang, 1998), in entorhinal cortex neurons (Alonso and Linhas, 1989; Dickson et al., 2000a), or in inferior olivary neurons (Linhas and Yarom, 1986)) pacemaking is clearly related to subthreshold membrane potential oscillations. However, these could not be isolated in WCRs by gradually hyperpolarizing Golgi cells with steady currents. This situation clearly resembles that observed in subthalamic neurons, in which pacemaking was extremely regular and subthreshold oscillations could not be decoupled from spikes (Bevan and Wilson, 1999). In the model, after switching-off spike-related mechanisms (i.e., $I_{\text{Na-t}}$, $I_{\text{K-V}}$, and $I_{\text{K-C}}$), subthreshold oscillations (30 mV peak-to-peak centered around threshold) emerged at nearly the same frequency (5 vs. 4.46 Hz) of spontaneous firing (Figure 9A). Several putative mechanisms might be responsible for these subthreshold oscillations in Golgi cells, exploiting the specific properties of the two depolarizing currents, $I_{\text{Na-p}}$ and $I_h$, and of the two
slow repolarizing currents, $k_{\text{slow}}$ and $k_{\text{AHP}}$. The role of each of these currents was considered in turn.

Blocking either $k_{\text{Na-p}}$ or $k_h$ modified the subthreshold oscillations although in quite a different way. By blocking $k_h$, the oscillation slowed down (from 4.46 to 3.3 Hz) and the original frequency could be restored by injecting a tonic depolarizing current (Figure 9A). Conversely, following 50% $k_{\text{Na-p}}$ reduction, the oscillation disappeared. In this case, a large tonic current injection compensated the partial block only to a very small degree (Figure 9B). Concerning the repolarizing feedback, blocking either $k_{\text{slow}}$ or $k_{\text{AHP}}$ increased oscillation amplitude but with opposite effects on frequency ($k_{\text{slow}}$ block reducing it to 2.9 Hz and $k_{\text{AHP}}$ block increasing it to 4.5 Hz), while blocking both $k_{\text{slow}}$ and $k_{\text{AHP}}$ led to a cessation of oscillation (Figure 9B).

The results reported in Figure 9A, B suggest that $k_{\text{Na-p}}$ and $k_h$ played distinct roles in Golgi cell subthreshold oscillations. The current plots in Figure 9C and the phase diagram in Figure 6D show that $k_{\text{Na-p}}$ underwent pronounced modulation during the cycle, while $k_h$ did not. Thus, while the brisk $k_{\text{Na-p}}$-activation was able to drive the cycle, the slow activation of $k_h$ provided a nearly tonic depolarization setting the cell into the appropriate membrane potential range. Concerning feedback repolarization, both $k_{\text{slow}}$ and $k_{\text{AHP}}$, which also showed pronounced dynamics during the cycle (Figure 9C, D), contributed to the feedback repolarization. The differential action of the two $K^+$ currents on cycle shape and frequency reflected their different voltage dependence and kinetics. It should be noted that the subthreshold oscillation reached potentials normally visited by spikes (about $-30 \text{ mV}$, cf. Figure 2A), so that $k_{\text{AHP}}$ could be activated causing a rapid and intense repolarization.

In accordance with the pharmacological tests reported in Figure 5 and the mechanisms reported in Figure 6, simulations shown in Figure 9 suggested that pacemaking was driven by subthreshold oscillations tightly coupled to spikes.

**DISCUSSION**

In this paper, we show that the theta-frequency pacemaking of cerebellar Golgi cells observed in vitro (Forti et al., 2006) can be explained by a computational model based on a mechanistic reconstruction of the underlying ionic processes. The model predicted that pacemaking depended on the actions of four ionic currents, $h$, $h_{\text{Na-p}}$, $k_{\text{AHP}}$, and $k_{\text{slow}}$ (Forti et al., 2006). $h$ brought membrane potential into the pacemaker region where the $h_{\text{Na-p}}/k_{\text{AHP}}/k_{\text{slow}}$ interaction generated pacemaking. Moreover, following hyperpolarization, $h$, $h_{\text{Ca-LVA}}$, and $k_{\text{Ca-LVA}}$ caused rebound excitation. In response to a depolarization, $k_{h}$ helped setting the delay for spike activation and $k_{h,T}$ (Alshahi et al., 2004; Magistretti et al., 2006) enhanced the initial doublet, whose separation from the rest of the discharge was sharpened by adaptation caused by $k_{\text{AHP}}$ and $k_{\text{slow}}$. Calcium-dependent regulation of ionic currents played also a critical role. By being coupled to $k_{h,C}$, $\Delta F_{\text{Ca-LVA}}$ enhanced the fast phase of spike AHP.
thereby resetting the spiking mechanism and sustaining high-frequency discharge. By being coupled to \( I_{K-AHP} \), \( I_{Ca-HVA} \) enhanced the slow phase of spike AHP and spike-frequency adaptation during repetitive discharge. The numerous currents revealed experimentally allowed the model to simultaneously control multiple aspects of electroresponsiveness, ranging from pacemaking to various response patterns elicited by depolarization and hyperpolarization. Each current could be attributed a characteristic role, as further explained in the following paper (Solinas et al., 2007), despite the apparent redundancy of some mechanisms (e.g., of \( I_{K-AHP} \) and \( I_{K-slow} \), which are both slow-repolarizing currents, or \( I_{Na-p} \) and \( I_{h} \), which could both sustain pacemaking). Moreover, the model allowed to define a typical, or “canonical,” Golgi cell electroresponsive pattern covering an extended experimental set of recordings (Forti et al., 2006). Apparently, histochemical differences (Geurts et al., 2001; Simat et al., 2007) did not remarkably impact on basic electroresponsiveness. Although the present model could be improved by more detailed analysis of channel gating and calcium dynamics in Golgi cells, it is a clear advancement over previously published models (Maex and De Schutter, 1998).

Pacemaking in the model was sustained by subthreshold oscillations caused by \( I_{Na-p} \) coupling with slow repolarizing currents, resembling mechanisms observed in some basal ganglia neurons (see Surmeier et al., 2005 for review). The \( I_{Ca-HVA} - I_{K-AHP} \) system influenced the oscillations but, unlike in a model of dopaminergic neurons of the substantia nigra (Wilson and Callaway, 2000), it was not essential for pacemaking.

Figure 6. Ionic currents and intracellular calcium changes. (A) The traces show the model response during alternated phases of pacemaking, depolarization, and hyperpolarization (the injected current is shown at the bottom). The voltage trace (Vm) shows the membrane potential change. The current traces (Im) show the ionic currents in the ISI. It should be noted that \( I_{Na-p} \), \( I_{K-AHP} \), and \( I_{K-slow} \), but not \( I_{h} \), show a marked voltage-dependent modulation during the pacemaker cycle. The calcium ([Ca\(^{2+}\)]\text{$_{i}$}) trace shows that [Ca\(^{2+}\)]\text{$_{i}$} is promptly reset to its background level of 50 nM following a transient increase during the spike. The inset shows the [Ca\(^{2+}\)]\text{$_{i}$} accumulation in coincidence with the high-frequency bursts and spike frequency adaptation. (B) The traces enlarge the model response during a single action potential (Vm). The ionic currents primarily responsible for action potential generation (Im) and the Ca\(^{2+}\) currents (\( I_{Ca-HVA} \)) are shown along with the [Ca\(^{2+}\)]\text{$_{i}$} transient. Note the delayed activation of the repolarizing mechanisms and of the slow time course of [Ca\(^{2+}\)]\text{$_{i}$} transient. \( I_{Na-p} \), \( I_{K-AHP} \), and \( I_{h} \) are much larger than currents governing the ISI (cf. panel A).

Figure 7. Ionic currents controlling response intensity and latency. (A) Effect of \( I_{Na-r} \) on the first two spikes elicited by a depolarizing current step. The model was hyperpolarized to −70 mV by tonic current injection and a step current injection of 600 pA elicited an initial firing rate of 200 Hz. \( I_{Na-r} \) block (middle trace) increased the first ISI (2 ms increase) reducing the initial firing rate to 142 Hz. \( I_{Na-r} \) did not affect the 1st spike latency but its activation after the 1st spike could accelerate membrane depolarization and generation of the 2nd spike. (B) The model was hyperpolarized to −70 mV by tonic current injection and then injected with a step current of 100 pA. The 1st spike latency was 18 ms and was reduced by \( I_{K-A} \) blockage to 12 ms. The lower traces show the stimulus protocols.
cell remains to be done, we point out that pharmacological block of SK channels generating $I_{\text{K-AHP}}$ did not arrest pacemaking in vitro (Forti et al., 2006). $I_h$ in our model proved critical to keep the oscillatory mechanisms in the appropriate membrane potential range for activation. This does not exclude that specific regulation of $I_h$ reversal potential might enable $I_h$ to take more actively part to oscillatory dynamics, as proposed for subthreshold oscillations in entorhinal neurons (Dickson et al., 2000a) and for pacemaking in cholinergic interneurons (Bennett et al., 2000). Finally, $I_h$ and $I_{\text{Ca-LVA}}$ did not give raise to any low-threshold pacemaking (Huguenard and McCormick, 1992; McCormick and Huguenard, 1992; McCormick and Pape, 1990) in Golgi cells but rather played a substantial role in rebound excitation, as also proposed for subthalamic neurons (Chan et al., 2004; Pennartz et al., 1997).

It may appear surprising that Golgi cell pacemaking was not associated with evident subthreshold oscillations, as it has been reported in, for example, amygdaloid neurons (Pape and Driesang, 1998), entorhinal cortex neurons (Alonso and Llinaís, 1989; Dickson et al., 2000a) or inferior olivary neurons (Llinás and Yarom, 1986). The model predicted that subthreshold oscillations in Golgi cells are tightly coupled to spikes, resembling the case of subthalamic neurons (Bevan and Pape, 1999, see also Shin et al., 2007). The regularizing action of $I_{\text{K-AHP}}$ occurs therefore through a repeated phase-reset of the pacemaker cycle, a concept fully developed in the continuation of this work (Solinas et al., 2007).

The intrinsic excitable properties of Golgi cells could have relevant implications for cerebellar signal processing. The basal discharge may become irregular in vivo due to synaptic noise suggesting that, similar to what was observed in a Purkinje cell model (De Schutter and Bower, 1994; Jaeger et al., 1997, see also Shin et al., 2007), in vivo Golgi cell spiking is still sustained by intrinsic pacemaking (see Solinas et al., 2007). The Golgi cell responses to protracted depolarizing pulses in vitro is characterized (after the initial burst) by a sustained discharge with a frequency depending on the intensity of the input. The modulation of response frequency would allow the Golgi cell to follow the temporal evolution of sensory-motor patterns (Miles et al., 1980). This confirms that the short-lasting Golgi cells spiking responses to long duration stimuli observed in vivo (Tahon et al., 2005) are due to depression of afferent input, as opposed to intrinsic mechanisms. Moreover, the presence of mechanisms enhancing the initial spike burst and regulating its initiation may be relevant for exerting precisely timed inhibitory effects on granule cells. An extended investigation of the effect of the ionic channel complement is presented in the companion paper (Solinas et al., 2007).
IK-AHP and IK-slow appear primarily involved in causing phase-reset and resonance, respectively. The model has therefore the ability to predict Golgi cell responses in various operating condition and appears as a useful tool for large-scale simulations of the cerebellar network.

CONFLICT OF INTEREST STATEMENT
The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

ACKNOWLEDGMENTS
This work was supported by grants CEREBELLUM (QLG3-CT-2001-02256) and SENSOPAC (FP6-IST028056) of the European Commission, by CNISM and Fondazione CARIPLO of Italy, by UA and FWO of Belgium, and by OIST of Japan. L.F. was the percipient of an INGENIO fellowship from the European Social Fund, Ministero del Lavoro e della Previdenza Sociale and Regione Lombardia of Italy.

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