A Non-Volatile Tunable Ultra-Compact Silicon Photonic Logic Gate

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Abstract: Logic gates, as one of the most important basic units in electronic integrated circuits (EICs), are also equally important in photonic integrated circuits (PICs). In this study, we proposed a non-volatile, ultra-compact all-photonics logic gate. The footprint is only 2 µm × 2 µm. We regulate the phase change of optical phase change materials(O-PCMs) Sb2Se3 to switch the function of the logic gate. The Sb2Se3 possesses a unique non-volatile optical phase change function; therefore, when Sb2Se3 is in the crystalline or amorphous state, our device can work as XOR gate or AND gate, and our designed logic ‘1’ and logic ‘0’ contrasts reach 11.8 dB and 5.7 dB at 1550 nm, respectively. Compared with other traditional optical logic gates, our device simultaneously has non-volatile characteristics, tunability, and additionally an ultra-small size. These results could fully meet the needs of fusion between PICs and EICs, and developing truly chip-scale optoelectronic logic solution.

Keywords: logic gate; phase change material; inverse design

1. Introduction

In recent years, large-scale photonic integrated systems have been confirmed to have a higher speed, higher capacity, and lower power consumption when performing computing tasks, which is suitable for the development of next-generation computing platforms [1,2]. In EICs, a field programmable gate array (FPGA) consists of a large number of logic units. Due to its flexible programmability and non-volatility, it has been widely used in electronics and communications. In PICs, a non-volatile tunable logic gate is also essential for implementing similar integrated devices. On the other hand, the all-optical logic gate is a key device for realizing optical communication networks, optical computing, and optical signal processing, it has been extensively studied in recent years. Many schemes have been proposed to realize all-optical logic gates, such as fibers [3–5], semiconductor optical amplifiers [6–8], photonic crystals [9–13], and Mach–Zehnder interferometers [14]. However, the functions of the logic gates implemented by these schemes are fixed and cannot be regulated.

The size of an optical logic gate based on the above schemes is generally tens to hundreds of square microns, being a significant fusion barrier between PICs and EICs. For example, in Ref. [15], the photonic crystal method was used to realize an all-optical logic gate with a size of 252 µm². In Ref. [16], also using the photonic crystal method, a gate with a size of 729 µm² was realized. In contrast, we used the inverse design...
method to realize a device with an ultra-small size; this method has been widely used in photonic devices design [17,18], such as in power splitters [19,20], focusing wavelength demultiplexers [21,22], and gratings [23,24]. Compared with traditional design methods, an inverse design method can search a larger parameter space to obtain high-performance and more compact photonic devices. To meet the needs of the development of photonic integrated circuits, many heuristic algorithms have been applied to an inverse design, including direct binary search algorithm (DBS) [25–28], genetic algorithm [29,30], particle swarm algorithm [31,32], objective-first algorithm [33,34], and neural networks [35–37].

O-PCMs are considered to be promising candidate material for designing reconfigurable photonic devices and memory units [38–40]. This is because their optical properties, such as the refractive index change drastically in a non-volatile manner. Devices based on O-PCMs have received signification research attention. They have been investigated for optical switches [41–44], metalens [45–47], mode converters [48,49], and optical neural networks [50,51]. In addition, O-PCMs have been used for memory cells due to their characteristics of fast phase transitions [52]. Experiments have shown [53] that the non-volatile phase transitions of O-PCMs between the crystalline and amorphous states can be controlled using light or electricity. Furthermore, Ge$_2$Sb$_2$Se$_5$ (GST) [54] is one of the widely used non-volatile O-PCMs. The complex refractive index of GST in the amorphous and crystalline states is 4.6-i0.12 and 7.45-i1.49 at 1550 nm, respectively. However, the extinction coefficient of GST crystalline state is too large, which could result in excessive loss. Sb$_2$Se$_3$ [55–57], as a new type of O-PCMs, it’s refractive index will change a lot during the transition between crystalline and amorphous states. More importantly, it has ultra-low loss in the commercial C-band. For example, at 1550 nm, the complex refractive index of Sb$_2$Se$_3$ in the amorphous and crystalline states is 3.285-i0 and 4.050-i0, respectively, where the imaginary parts are zero for both states, indicating ignorable light absorption. Both crystalline and amorphous states show ultra-low loss, being a rare feature among many O-PCMs.

In this study, we designed and demonstrated a non-volatile, tunable, ultra-compact optical logic gate on a silicon-on-insulator (SOI) platform based on the DBS algorithm and Sb$_2$Se$_3$. The footprint of the device is only 2 µm × 2 µm. For our device, when Sb$_2$Se$_3$ is in the crystalline state, the XOR gate is realized, and when Sb$_2$Se$_3$ is in the amorphous state, it can be used as an AND gate. The function of the device can be regulated by switching the state of Sb$_2$Se$_3$. As far as we know, most existing optical logic gates have a large footprint, which is not conducive for integration, and their function cannot be controlled. Here, we first propose a tunable ultra-compact silicon optical logic gate, which can be used as a logic gate in future photonic integrated circuits.

2. Chip Design and Algorithm

The 3D structure, the size of the proposed logic gate and the complex refractive index and atomic distribution of Sb$_2$Se$_3$ are shown in Figure 1. Our device was based on the SOI platform, where the buried layer was 3 µm silicon dioxide and the top layer was standard 220 nm silicon. Furthermore, Sb$_2$Se$_3$ was embedded in the top layer silicon. Here, we used silicon instead of Si$_3$N$_4$ or other materials, mainly considering the complete compatibility with CMOS processes. Moreover, silicon has about a 3.5 refractive index at around 1550 nm and then shows a very strong light field confinement. Before using the DBS algorithm, the 2 µm × 2 µm design area was divided into 20 × 20 square pixels for digital binarization, the size of each pixel was 100 nm × 100 nm. Each pixel has two states, ’0’ and ’1’, where ‘0’ represents silicon, and ’1’ represents Sb$_2$Se$_3$. The silicon at the ’1’ positions was etched to a depth of 220 nm and replaced by Sb$_2$Se$_3$. Compared with the O-PCMs on the surface of silicon, embedded O-PCMs in silicon could have a stronger ability to control light field. With design experience, we set the width of the two input waveguides as 400 nm, the distance between them was 800 nm, and the width of the output waveguide as 400 nm. In Figure 1d, the atomic distribution in the crystalline and amorphous states of Sb$_2$Se$_3$ and its complex refractive index at 1500–1600 nm are presented. Especially, for Sb$_2$Se$_3$,
extinction coefficient (k) is 0 in both crystalline and amorphous state. The design goal of our device is that the device should act as an XOR gate when Sb$_2$Se$_3$ is in the crystalline state and an AND gate when Sb$_2$Se$_3$ is in the amorphous state.

![Figure 1](image_url)

**Figure 1.** Schematic of the logic gate device. (a) Top view of the device; (b) Side view of the device; (c) Three-dimensional diagram of the device showing two input ports in the left and one output port in the right; (d) Atomic distribution and complex refractive index of Sb$_2$Se$_3$ in crystalline and amorphous states.

The flowchart of the inverse design based on the DBS algorithm is shown in Figure 2. The structure of the design area was represented by a 20 × 20 matrix of 0 s and 1 s. The corresponding spectrum of each structure was obtained after a simulation. We defined a figure of merit (FOM) to evaluate the performance of the current structure. In this way, we abstracted the physical optimization objective into a mathematical optimization process. With the DBS algorithm, we optimized the structure and improved the FOM to obtain the required device performance. FOM was defined as follows:

\[
FOM_C = (P_{C10} + P_{C01}) - 3 \times P_{C11} - |P_{C10} - P_{C01}| \quad (1)
\]

\[
FOM_A = P_{A11} - 1.5 \times (P_{A10} + P_{A01}) - 3 \times |P_{A10} - P_{A01}| \quad (2)
\]

\[
FOM = 1.5 \times FOM_C + FOM_A \quad (3)
\]

The FOM consisted of two parts, FOM$_A$ and FOM$_C$, as a crystalline and amorphous FOM, respectively. Each source was assumed to have four states, i.e., '00', '01', '10' and '11'; the first represented the state of input1 and the second represented the state of input2. Furthermore, $P_{C10}$, $P_{C01}$, and $P_{C11}$ in formula (1) are the output intensities when the input states are '00', '01' and '11', respectively, in the crystalline state, and $P_{A10}$, $P_{A01}$, and $P_{A11}$ in formula (2) are amorphous state. During the optimization process, our FOM is adjusted to obtain the final form. A 20 × 20 matrix containing 0 s and 1 s was first randomly generated. Then, the 3D FDTD was used to solve Maxwell’s equations and the FOM of the device structure was calculated. A starting point in the matrix was selected and its state was flipped ('0' to '1' or '1' to '0'), and the FOM was calculated after the flip. If the FOM improved, the flipped state was retained, and if the FOM did not improve, the flipped state was restored. This procedure was repeated for each point in the structure (by row
or column). After all the points in the matrix are calculated, one iteration ends and a new iteration starts again. The algorithm will continue to run until the target conditions are met.

3. Chip Simulations and Analysis

All simulations in this work were conducted with the 3D FDTD analysis software (Ansys Lumerical FDTD 2020 R2.4). In the case that our proposed device functions as an XOR gate, when the source is ‘00’ or ‘11’, the output state is ‘0’, and when the source is ‘10’ or ‘01’, the output state is ‘1’; In cases where the device functions as an AND gate, when the source is ‘10’, ‘01’, or ‘11’, the output state is ‘0’, and when the source is ‘11’, the output state is ‘1’.

We used the contrast ratio (CR) to evaluate the degree of difference between the ‘0’ and ‘1’ states of the logic gate device, where CR is defined as:

$$CR = 10 \times \log \left( \frac{P_1}{P_0} \right)$$

(4)

where $P_1$ is the minimum intensity with a logical value of ‘1’, and $P_0$ is the maximum intensity with a logical value of ‘0’.

When the source state is ‘00’, the output state of our device is always ‘0’; hence, only the three states of ‘10’, ‘01’ and ‘11’ are discussed in the following section. Figure 3 presents the optimized structure of our device and the energy density distribution of the three states of ‘10’, ‘01’, and ‘11’ when our device functions as the XOR gate. As shown in Figure 3c,d, when the input state is ‘01’ or ‘10’, the power of the output port is higher, which is the...
‘1’ state; when the input state is ‘11’, almost no output power is observed, which is the ‘0’ state. From Figure 3b, it can be seen that after the light of inputs 1 and 2 enters the device, it converges above the output port and cannot escape through the output port, making the output power low. Figure 4 presents the spectral power curve of the XOR gate when Sb$_2$Se$_3$ is in the crystalline state at 1530 to 1560 nm.

Table 1 shows the output intensity of the XOR gate in each source state when Sb$_2$Se$_3$ is in the amorphous state and our device functions as an AND gate. Figure 5 shows the energy density distribution of the states of ‘10’, ‘01’, and ‘11’ when Sb$_2$Se$_3$ is in the crystalline state at 1530 to 1560 nm.

Table 1. Truth table for the all-optical XOR logic gate.

| Input 1 | Input 2 | Output | Threshold | Output State | XOR Gate Output |
|---------|---------|--------|-----------|--------------|----------------|
| 0       | 0       | 0      | 0.2 P$_{in}$ | 0            | 0              |
| P$_{in}$| 0       | 0.381 P$_{in}$ | 0.2 P$_{in}$ | 1            | 1              |
| 0       | P$_{in}$| 0.381 P$_{in}$ | 0.2 P$_{in}$ | 1            | 1              |
| P$_{in}$| P$_{in}$| 0.025 P$_{in}$ | 0.2 P$_{in}$ | 0            | 0              |
Figure 5 shows the energy density distribution of the states of ‘10’, ‘01’, and ‘11’ when Sb$_2$Se$_3$ is in the amorphous state and our device functions as an AND gate. Figure 5c,d shows the energy distribution when the input state is ‘01’ and ‘10’, respectively. The output power is relatively low and the output is the ‘0’ state. When state is ‘11’, as shown in Figure 5b, the output power is high and the output is the ‘1’ state. Figure 6 is the spectral power curve of the AND gate when Sb$_2$Se$_3$ is in the amorphous state in the range from 1530 to 1560 nm.

Table 2 shows the output intensity of the AND gate in each source state when λ = 1550 nm. When the source state is ‘00’, the output intensity is 0 and the output state is logic ‘0’. For the source states of ‘10’, ‘01’, and ‘11’, the output intensities and states are 0.213 $P_{in}$ and ‘1’, 0.213 $P_{in}$ and ‘1’, and 0.783 $P_{in}$ and ‘0’, respectively. The CR of the AND gate is 5.7 dB, and we set the threshold between the logics ‘0’ and ‘1’ to 0.5 $P_{in}$. The above results show the performance parameters of our device as the XOR and the AND gates. In addition, it is worth mentioning that due to the ultra-small size of our device, its response time is less than 1 ps. These properties enable its use in photonic integrated circuits.

Figure 5. Energy density distribution of the AND gate when Sb$_2$Se$_3$ is in the amorphous state. (a) The final structure of our device; the red pixels are Sb$_2$Se$_3$ in the amorphous state and the purple pixels are silicon; (b) Energy distribution when the input state is ‘11’; (c) Input state is ‘01’; (d) Input state is ‘10’.

Figure 6. Performance graph of AND gate. (a) Spectral power curve of the AND gate when Sb$_2$Se$_3$ is in the amorphous state; (b) Contrast ratio curve of the AND gate.
Table 2. Truth table for the all-optical AND logic gate.

| Sb$_2$Se$_3$ is in the Amorphous State and Functions as an AND Gate |
|---------------------------------------------------------------|
| Input1 | Input2 | Output | Threshold | Output Logic State | AND Gate Output |
|--------|--------|--------|-----------|--------------------|-----------------|
| 0      | 0      | 0      | 0.5 $P_{in}$ | 0                  | 0               |
| $P_{in}$ | 0      | 0.213 $P_{in}$ | 0.5 $P_{in}$ | 0                  | 0               |
| 0      | $P_{in}$ | 0.213 $P_{in}$ | 0.5 $P_{in}$ | 0                  | 0               |
| $P_{in}$ | $P_{in}$ | 0.783 $P_{in}$ | 0.5 $P_{in}$ | 1                  | 1               |

Here, we discuss the fabrication errors and tolerance performance of our device. In actual device manufacturing, it is common that over-etching or under-etching causes errors in device manufacturing that affects device performance [58]. For our device, the etching depth should be 220 nm; however, in actual manufacturing, the etching is very likely to be insufficient, not reaching 220 nm. Therefore, we analyzed the performance of the device in the case of insufficient etching. We assumed that the depth of each etched Sb$_2$Se$_3$ hole was in a random range between 180–220 nm, and assumed E1, E2, E3, and E4 to be devices with errors manufactured according to this rule. Figure 7 shows the results. Figure 7a–c plot the numerical results of the XOR gate and source states of ‘01’, ‘10’, ‘11’, and the results of the AND gate are plotted in Figure 7d–f. It can be seen from Figure 7 that for the devices E1, E2, E3, and E4 with manufacturing errors, although the etching depth of each pixel is a random value between 180–220 nm, their spectrum did not change significantly, and the functionality of our device is still realized. In our standard device, the CR of the XOR gate is 11.8 dB, the CR of AND gate is 5.7 dB, as for E1, E2, E3, and E4, the lowest CR of the XOR gate is 6.9 dB and that of the AND gate is 5.1 dB. In this comparison, although the CR of the XOR gate decreased significantly, the threshold between the designed logic states of ‘0’ and ‘1’ is 0.2 $P_{in}$, the intensity of logic ‘0’ is lower than 0.2 $P_{in}$, and that of logic ‘1’ is higher than 0.2 $P_{in}$. The four under-etched devices still realize the XOR gate function, as for the AND gate, the error is very small and its functioning remains virtually unchanged. The above analysis shows that our device can realize the corresponding function even if the etching depth is not up to the requirement during manufacturing, indicating that the device is robust.

Logic gates can be combined with some common photonic devices to realize more complex devices. Here, we show two possible integration methods to realize the ultra-compact photonics half adder and full adder. As shown in Figure 8a, our logic gates, beam splitter and waveguide crossing can be combined to form a half adder. Figure 8b is a full adder, which can be composed of a half adder, OR gate and bend waveguide. Furthermore, by combining the full adders, a complete addition operation can be realized. In addition, the combinational logic circuits, encoders, and decoders similar to EICs can be implemented by logic gate devices combined with common photonic devices.

Furthermore, we simulated the performance of the ultra-compact photonics half adder and full adder and provide examples in Figure 9. Additionally, Figure 9a presents the example of a logic timing diagram of the half adder, where the input states are ‘00’, ‘10’, ‘11’, ‘01’, ‘01’, ‘11’, presented by 2.5 picosecond time interval pulses in sequence. The Sum output was logical ‘0’, ‘1’, ‘0’, ‘1’, ‘0’, ‘1’, ‘0’, respectively, and the Carry output was ‘0’, ‘0’, ‘0’, ‘1’, ‘0’, ‘0’, ‘1’. The full logic truth table of half adder is shown in Table 3. By comparing Table 3 and Figure 9a, we found that the photonics half adder can implement the logic functions, with a contrast ratio of the ‘1’ and ‘0’ of the Sum channel reaching 6.9 dB, and the Carry channel being 5 dB. Next, Figure 9b shows the results of full adder logic gates. In order to keep the input intensity of the half adder consistent, we set the input intensity of Cin equal to the logical ‘1’ Sum output intensity of the half adder. The state of Cin, Input1 and Input2 are represented by the first bit, the second bit and the third bit, respectively. The input pulse sequence is ‘111’, ‘000’, ‘001’, ‘010’, ‘100’, ‘011’, 2.5 picosecond time per state, while the Sum output was logical ‘1’, ‘0’, ‘1’, ‘1’, ‘1’, ‘0’, respectively, and the Carry output was ‘1’,
'0', '0', '0', '0', '1', which is identical to the truth in in Table 4. As a result, the contrast ratio of the '1' and '0' of the Sum channel is about 6.9 dB, and the Carry channel is 5 dB, being similar with a half adder. Additionally, during our FDTD calculation, the example pulse interval is 2.5 picoseconds. That is, the corresponding data sequence in the photonics logic gate is 400 Giga bits per second. Moreover, if considering the extremely small size of the device, our device may support femtosecond temporal pulses, or dozens of Tera bits per second for ultra-fast data sequence.

Figure 7. Analysis of fabrication tolerances. Comparison of the intensity between the standard device and devices with an etching error. (a) XOR gate with the input state of '01'; (b) XOR gate with the input state of '10'; (c) XOR gate with the input state of '11'; (d) AND gate with the input state of '01'; (e) AND gate with the input state of '10'; (f) AND gate with the input state of '11'.

Figure 7. Analysis of fabrication tolerances. Comparison of the intensity between the standard device and devices with an etching error. (a) XOR gate with the input state of '01'; (b) XOR gate with the input state of '10'; (c) XOR gate with the input state of '11'; (d) AND gate with the input state of '01'; (e) AND gate with the input state of '10'; (f) AND gate with the input state of '11'.

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Figure 8. Schematic diagram of half adder and full adder composed of logic gates. (a) Half adder composed by two beam splitters, a waveguide crossing, a XOR gate and a AND gate; (b) Full adder composed by two half adder and a OR gate.

Figure 9. The performance of the half adder and full adder. (a) Input pulse and output intensity of half adder; (b) Input pulse and output intensity of full adder.
Table 3. Logic truth table of the half adder.

| Input1 | 0 | 1 | 0 | 1 |
|--------|---|---|---|---|
| Input2 | 0 | 0 | 1 | 1 |
| Sum    | 0 | 1 | 1 | 0 |
| Carry  | 0 | 0 | 0 | 1 |

Table 4. Logic truth table of the full adder.

| Cin | 0 | 0 | 0 | 1 | 0 | 1 | 1 | 1 |
|-----|---|---|---|---|---|---|---|---|
| Input1 | 0 | 0 | 1 | 0 | 1 | 0 | 1 | 1 |
| Input2 | 0 | 1 | 0 | 0 | 1 | 1 | 0 | 1 |
| Sum | 0 | 1 | 1 | 1 | 0 | 0 | 0 | 1 |
| Carry | 0 | 0 | 0 | 0 | 1 | 1 | 1 | 1 |

4. Conclusions

In this work, we used the DBS algorithm, combined with O-PCM Sb$_2$Se$_3$ to design a tunable, ultra-compact silicon photonic logic gate. At 1550 nm, when Sb$_2$Se$_3$ is in the crystalline state, our device functions as an XOR gate, and its CR is 11.8 dB. When Sb$_2$Se$_3$ is in the amorphous state, it can work as an AND gate and has a CR of 5.7 dB. The footprint of our device is only 2 $\mu$m $\times$ 2 $\mu$m. We also studied the manufacturing tolerances for device etch depths ranging from 180–220 nm. With thresholds of the ‘1’ and ‘0’ states of our XOR and AND gates of 0.5 $P_{in}$ and 0.2 $P_{in}$, respectively, the device can still achieve corresponding functions. Furthermore, we provide a possible half adder and full adder structure, which can serve as a kind of optoelectronic fusion module. The inverse design method combined with O-PCM has great potential in the design of ultra-compact and tunable devices. We believe that such logic gates could have various application prospects in future ultra-high density, reconfigurable, scalable photonic integrated circuits.

Author Contributions: Conceptualization, Z.P. and J.F.; methodology, Z.P., H.L. and J.F.; software, Z.P., H.Y. and J.F.; writing original draft preparation, Z.P., Y.W., S.Z. and W.C.; resources, J.Y. and J.W.; visualization, X.R., Y.S. and H.C.; supervision, J.Y. and J.W.; funding acquisition, J.Y. and J.W. All authors have read and agreed to the published version of the manuscript.

Funding: This work is funded by the National Natural Science Foundation of China (60907003, 61805278, 61875168); Chongqing Science Funds for Distinguished Young Scientists (cstc2021jcyj-jqX0027); China Postdoctoral Science Foundation (2018M633704, 2017M612885); Foundation of NUDT (JC13-02-13, ZK17-03-01); Hunan Provincial Natural Science Foundation of China (13JJ3001); Chongqing Postdoctoral Science Foundation Special Funded project (Xm2017008); Program for New Century Excellent Talents in University (NCET-12-0142); the Fundamental Research Funds for the Central Universities (No. XDJK2013C117); Ph.D. Fund of Southwest University (No. 20130553); Guangxi Key Laboratory of Cryptography and Information Security (GCS201801).

Data Availability Statement: Data available in a publicly accessible repository.

Conflicts of Interest: The authors declare no conflict of interest.

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