A Baseline Correction Algorithm for FAST

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Abstract

The Five-hundred-meter Aperture Spherical radio Telescope (FAST) is the most sensitive ground-based, single-dish radio telescope on Earth. However, the original HI spectra produced by FAST are affected by standing waves. To maximize the power of FAST for high-sensitivity observations, we proposed an algorithm that combines fast Fourier transforms and extreme envelope curves to automatically correct the baselines of FAST HI spectra and remove standing waves from the baselines. This algorithm can reduce the amplified noise level caused by standing waves to a near-ideal level without losing signals or introducing false signals. The root mean square of the average baseline reaches \(~8\) mK, approaching the theoretical sensitivity of an HI spectrum produced by FAST for an integration time of 335 minutes, i.e., \(~6\) mK.

Key words: methods: observational – techniques: spectroscopic – telescopes

1. Introduction

The Five-hundred-meter Aperture Spherical radio Telescope (FAST), located in Guizhou Province of Southwest China, is the world’s largest single-dish radio telescope. It is an important facility for surveying neutral hydrogen up to the edge of the universe, detecting weak space signals, hearing possible signals from other civilizations, etc. (Nan et al. 2011; Qian et al. 2020). Although it has produced significant scientific achievements (Qian et al. 2020), the baselines of the original HI spectra are not flat enough and they contain massive standing waves that might be generated by reflections between the dish and the receiver cabin (Jiang et al. 2020). Although efforts have been made to minimize the standing wave effects of FAST data (Jiang et al. 2020), standing waves still exist in the spectra. Standing waves can amplify the noise level of a signal; e.g., ripple amplitudes of \(~15\) mK are commonly seen (Jiang et al. 2020), which cause the noise of the obtained high-sensitivity spectra to be higher than the theoretically predicted noise. Some studies usually require extreme high-sensitivity observations during the analysis of the spectra; e.g., searching for the stellar winds, compact high-velocity clouds and active star-forming dwarf galaxies (Lizano et al. 1988; Giovanardi et al. 1992; Burton et al. 2001; Salzer et al. 2002; Li et al. 2022). A high-precision baseline correction method is hence necessary for studies that require extreme high-sensitivity spectra.

Polynomial fitting and trigonometric function fitting methods are commonly used to make baseline corrections (Gan et al. 2006; Baek et al. 2011). These methods usually need to cut peaks from the original spectrum and estimate the baseline using a polynomial or trigonometric function. However, these methods can be ineffective if the baseline is complex or the format of the function is not good enough. The asymmetrically reweighted penalized least squares (arPLS) algorithm, developed from penalized least squares methods (Eilers 2003; Carlos Cobas et al. 2006; Zhang et al. 2010; Baek et al. 2015), is a widely used baseline correction method for FAST data (Wang et al. 2022; Zhang et al. 2022). The baseline can be estimated by changing the “weight” parameter iteratively. Similar weights are assigned to baseline regions without peaks, while no weights or small weights are assigned for peaks; once assigned, the weights gradually reduce as the level of the signal increases. However, arPLS is not good at removing standing waves.

To correct an inclined baseline and remove standing waves from the original HI spectra automatically produced by FAST, an algorithm combining fast Fourier transforms (FFTs) and extreme envelope curves (EECs), called FFTEEC, is proposed in this work.

2. Data

2.1. Observations

FAST is equipped with a 19 beam receiver and dual linear polarizer (i.e., XX and YY). The full bandwidth of the L-band is 500 MHz over the frequency range 1.0–1.5 GHz. The frequency resolution of the high-resolution modes is 476.84 Hz, corresponding to a velocity resolution of \(~0.1\) km s\(^{-1}\) @ 1.4 GHz. The beam size is 2.9', and the pointing error is \(~0.2\) (Li et al. 2018; Jiang et al. 2019, 2020). The data used in this paper consist of HI spectra of G176.51+00.20 observed on 2021 August 19 and 20, with the 19 beam tracking observing mode, and the total integration time of the data was 335 minutes with a sampling rate of one second. We resampled the data to a velocity resolution of 0.1 km s\(^{-1}\), and only considered data in the velocity range...
−7000–7000 km s$^{-1}$, i.e., 140,001 channels. The data displayed in this paper are from the $YY$ polarization of Beam M01.

### 2.2. Baseline of FAST

Figure 1 displays the original spectra after correcting the flux and velocity (Jiang et al. 2020). In the left panel, the non-uniform waterfall image indicates that the baselines of the different original spectra are inconsistent, and the average spectrum with massive standing waves is inclined. The right panel shows fringes in the waterfall image after roughly correcting the inclined baseline with the polynomial fitting method. The standing waves are obvious and unstable (see the partially enlarged view). The arc-shaped standing waves in the 2D waterfall image mean that the phases of the standing waves drift with time and the periods of standing waves in different spectra are inconsistent at the same time.

### 3. Method

Due to the irregular phases and periods of standing waves, as well as huge amounts of data produced by FAST, e.g., 6.0 TB for our data, we propose a highly precise baseline estimation algorithm, FFTEEC, to correct baselines automatically. First, the original spectrum is preprocessed by a polynomial fitting method. Second, the standing waves are extracted and removed using FFTs. Third, the EEC method is employed to calibrate the unsmooth parts in the baseline. Fourth, the extracted signal is combined with the baseline obtained in the third step as the result. Figure 2 displays the whole pipeline of the algorithm.

#### 3.1. FFTs

An FFT is a fast algorithm that computes discrete Fourier transforms (DFTs, Cooley & Tukey 1965). The DFT of sequence $x$ with length $N$ can be expressed as

$$X(k) = \text{DFT}[x(n)] = \sum_{n=1}^{N} x(n) \exp\left[-j2\pi nk/N\right], \quad 1 \leq k \leq N.$$ (1)

Any frequency with a large amplitude in the frequency domain, $X$, can be considered as a standing wave. A new sequence, $Y$, is obtained by sorting $X$ from high to low and removing the first several items.

#### 3.2. EECs

An EEC, $y_e$, can be obtained with the following steps. First, get the smoothed sequence $y$ from $Y$, and the local maximum and minimum are extracted from $y$. Second, the maximum and minimum envelope curve of $y$, $y_{\text{max}}$ and $y_{\text{min}}$, respectively, can be obtained by fitting the local extrema with a linear interpolation method. Thus, $y_e$ can be expressed as

$$y_e = \frac{y_{\text{max}} + y_{\text{min}}}{2}.$$ (2)

#### 3.3. Automatic Signal Extraction

To extract a complete signal, we propose an iterative method to cut the signal automatically. The initial signal range is obtained from the EEC, where $y_e$ is greater than $20\sigma$ of the smoothed baseline. The signal range needs to be extended if the difference between the wing of the smoothed signal and the baseline is...
greater than 1σ until it is smaller than 1σ. The method is displayed in Figure 2 as a subprocess that is framed by a box.

4. Performance of FFTEEC

4.1. Simulation

The simulated data considered here contain the pure analytical signal p, a simulated baseline, b, and random noise, j, which can be expressed as

\[ s(n) = p(n) + b(n) + 0.1 \times j(n), \]  

where p is

\[ p(n) = 100 \exp\left( -\frac{n^2}{25^2}\right) . \]  

and b contains the residuals extracted from the real data by the method in this paper.

Figure 2. Pipeline of FFTEEC. The figure corresponding to each step is the average spectrum after the step is processed.

Figure 3. Left: Results processed by different methods. The blue line shows the FFTEEC result, where the zero-point is 1, and the green line is the arPLS result. The flat baselines indicate that both methods can be used to correct inclined baselines. Right: Simulated standing waves and residuals from different methods. The blue line is the FFTEEC result and the green line is the arPLS result. The orange line is the simulated standing wave. The spectra have been separated by 1 on the y-axis to aid presentation. The profiles of the residuals of arPLS and FFTEEC are consistent with the simulated baseline. Standing waves are obvious in the residuals of FFTEEC, but unclear in the residuals of arPLS since the random noise is very high.

3 The random noise was generated by the random number generator numpy.random.randn from the Python language.

4 We used the C++ arPLS software package provided by Gabriel Kronberger to speed up the calculations. https://github.com/heal-research/arPLS.
in the bottom panel. The signals in the spectra processed by FFTEEC are consistent with the simulated signal, and the residuals between them are flat. In the simulated data, the width of the 3σ signal is 60.2, while that obtained by the automatic signal extraction method is 84.8, which is wider than the signal.

The root mean square error (RMSE) can be applied to illustrate the difference between two spectra, which can be expressed as

$$RMSE = \sqrt{\frac{1}{N} \sum_{n=1}^{N} (y(n) - \bar{y}(n))^2},$$

where $N$ is the length of the sequence, and $y$ and $\bar{y}$ are the signals to be compared. The RMSE between the simulated signals and the spectra processed by FFTEEC is 0.10, and that between the simulated signals and the simulated signals after random noise injection is also 0.10, indicating that there is only random noise in our result. As highlighted by the visual effect shown in Figure 4 and RMSE values of the simulated data, our method does not lose signals or introduce false signals.

The root mean square (rms) can also be employed to judge the effectiveness of our method, which can be expressed as

$$rms = \sqrt{\frac{1}{N} \sum_{n=1}^{N} (y(n))^2}.$$  

where $N$ is the length of the sequence and $y$ is the baseline after removing the signal. We compared the rms noise levels between the random noise of simulated data and our results, and the ratio between them is 1.12, indicating that our method can get a near-ideal rms noise levels.

4.2. Application to Real Data

We applied a 10th order polynomial fitting method to correct the inclined baseline preliminarily. However, the baseline is roughly flat after polynomial correction, but still bumpy, and there are massive standing waves with different periods in it.

Figure 5 presents the baseline of FAST in the frequency domain processed by FFT. The blue line is the original data, and there are three distinct standing waves with different frequencies. The orange line is the result where the standing waves have been removed. Here, we find that the frequencies of the standing waves of most FAST data are relatively consistent and contain three different frequencies, as shown in the figure after processing the data from different beams. However, some standing waves appear to have additional frequencies; e.g., the XX polarization of Beams M10, M13 and M15, the YY polarization of Beams M06, M07 and M17, etc. Taking the first 20 orders can remove almost all standing waves in the different data sets, whether the standing waves contain three frequencies or more. Thus the order of the FFT algorithm usually can be assigned to 20, indicating removal of the top 20 items with the largest amplitudes in the frequency domain. This number was obtained after experimenting using various FAST data sets.

Figure 6 displays spectra corrected by FFTEEC and arPLS, in the left and right panels respectively. In the left panel, the standing waves have been removed from the spectrum, since there is only random noise in the waterfall image. Meanwhile, FFTEEC is stable for the different spectra, since the average spectrum is so flat that it is hard to find any obvious standing waves. The result processed by arPLS is shown in the right panel; no matter in the waterfall image or average spectrum, standing waves are obvious, indicating that arPLS can just correct the inclined baseline and is not good at removing standing waves in the original H I spectra produced by FAST.
Table 1 lists the rms noise of the average spectrum of five groups, i.e., the XX polarization of Beams M02 and M03 and the YY polarization of Beams M01, M06 and M13. The first row is for arPLS and the second row is for FFTEEC. The parameter $\lambda$, which is used to control the balance between fitness and smoothness, of arPLS is $10^{11}$ when real data were processed. Based on the rms estimation method of the average spectrum provided by FAST (see Equation (10) of Jiang et al. 2020), the theoretical rms of the average spectrum, integrated for 335 minutes, is $6 \sim 8$ mK. The rms noise of arPLS is $\sim 20$ mK, which is significantly higher than theoretical rms if the standing waves have not been removed.

In conclusion, although arPLS can be applied to correct the inclined baseline, it is hard to remove standing waves from the baseline. In contrast, FFTEEC can effectively remove standing waves and obtain a near-ideal rms. There are some shortcomings for FFTEEC, e.g., it can only be used to correct the baseline when knowing the position of the signal, but it cannot automatically extract the signal as arPLS can. Additionally, sometimes the parameters of FFTEEC affect the results; e.g., the length of the smoothing box when automatically extracting a signal.

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