Towards Visible Light Hydrogen Generation: Quantum Dot-Sensitization via Efficient Light Harvesting of Hybrid-TiO$_2$

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We report pronounced enhancement of photoelectrochemical hydrogen generation of a quantum dot-sensitized hybrid-TiO$_2$ (QD/H-TiO$_2$) electrode that is composed of a mesoporous TiO$_2$ layer sandwiched by a double sided energy harvesting layer consisting of a surface-textured TiO$_2$ inverse opals layer on the bottom and a patterned mesoporous TiO$_2$ layer on the top. CdSe/H-TiO$_2$ exhibits a maximum photocurrent density of $\sim 16.2$ mA/cm$^2$, which is 35% higher than that of the optimized control sample (CdSe/P25), achieved by matching of the bandgap of quantum dot-sensitization with the wavelength where light harvesting of H-TiO$_2$ is observed. Furthermore, CdSe/H-TiO$_2$ under filtered exposure conditions recorded current density of $\sim 14.2$ mA/cm$^2$, the greatest value in the visible range. The excellent performance of the quantum dot-sensitized H-TiO$_2$ suggests that alteration of the photoelectrodes to suitable nanostructures with excellent light absorption may offer optimal strategies for attaining maximum efficiency in a variety of photoconversion systems.
The photoelectrode is thus needed and careful design of quantum dot-sensitized devices based on a light harvesting mechanism is crucial to maximize the efficiency in photocatalysis.

Here, we report an optimized quantum dot-sensitized hybrid-TiO$_2$ (QD/H-TiO$_2$) electrode sandwiched with light harvesting layers. The H-TiO$_2$ electrode is constructed using a 4.5 µm-thick mesoporous TiO$_2$ film made of crystalline nanoparticles (25 nm) and additionally a 500 nm-thick patterned mesoporous TiO$_2$ layer on the top and 1 µm-thick surface-textured TIO (st-TIO) layer on the bottom, both with diameters on the order of hundreds of nanometers, with the capability of AR and photon trapping, respectively, contributing to maximized light harvesting. The configuration of the light harvesting layer sitting on both the bottom and top of the mesoporous TiO$_2$ layer provides devices with similar mechanism to mirrors in a laser, where entrapped light is enhanced by repeated reflection at both ends and thus maximizes the intensity of the absorbed light at a wider solar spectrum. Furthermore, CdSe/H-TiO$_2$ greatly broadens the range of photocatalysis below the band gap of H-TiO$_2$ by enhanced quantum dot-sensitized electron transport substantially supported by the light harvesting mechanism of H-TiO$_2$. Under AM 1.5 G simulated sunlight illumination, the maximum photocurrent density of CdSe/H-TiO$_2$ reaches ~16.2 mA/cm$^2$, which is 33% higher than that of the optimized control sample (CdSe/P25), at an applied bias of 0.5 V versus Ag/AgCl. More surprisingly, the photocurrent from hydrogen generation of CdSe/H-TiO$_2$ obtained under UV-visible irradiation. The greatly improved performance of CdSe/H-TiO$_2$, which is attributable to the ability of light harvesting of H-TiO$_2$, suggests the strong need for engineering light in photocatalysis systems.

**Results**

H-TiO$_2$ was fabricated by the following four steps, as shown in Figure 1a. First, surface-activated polystyrene beads (PS) obtained by treatment with Triton X-100 were self-assembled using Colvin’s method into a few layers of a hexagonal structure on a FTO substrate with a thickness of 1 µm.

PS structures were infiltrated with a TiCl$_4$ aqueous solution and annealed in an air condition, inducing gelation of a TiO$_2$ sol and calcination of PS simultaneously, and thus forming the TIO structure (I). Second, a surface textured titanium inverse opal (st-TIO, II) coated with a diameter of 7 nm TiO$_2$ mesoporous structures was created by selective removal of triblock copolymer domains from the coated layer containing TiO$_2$ precursors and block copolymers on TIO. In principle, the st-TIO has been designed to minimize the loss of surface area for reaction sites while achieving a photon trapping effect. Third, P25 mesoporous structures (III) were deposited onto the st-TIO with a thickness of 5 µm. Fourth, a top-patterned mesoporous P25 nanostructure (p-P25, IV) was created by impressing using a 500 nm-thick PDMS mold with 500 nm diameter and 600 nm pitch square patterns followed by the sintering process. The targeted structure is H-TiO$_2$, composed of a 4.5 µm-thick mesoporous TiO$_2$ layer with a double sided energy harvesting layer comprising a 1 µm-thick st-TIO layer of a few hundred nanometer scale on the bottom and a 500 nm thick mesoporous TiO$_2$ square pattern on the top, giving a total thickness of 6 µm. Figure 1b is a scanning electron microscope (SEM) image of the st-TIO structure (II) with air holes with a diameter of 290 nm. The dimension of the TIO was determined so as to efficiently scatter light for maximized light trapping as optimized in our previous reports. Figure 1c is a SEM image of patterned mesoporous TiO$_2$ structures obtained by stamping a P25 layer using the PDMS molds followed by final sintering. The dimensions were fixed for the maximized AR effects, based on the previous reports. The volume fraction of TiO$_2$ in the AR layer is found to be 0.46 based on SEM imaging and the effective refractive index of the AR layer is 1.62 as calculated in the Supporting Information. Figure 1d is a cross-sectional image of the H-TiO$_2$ photoanode (V) sandwiched by a thin upper-layer of AR and a photon trapping under-layer. The surface area of the 6 µm-thick H-TiO$_2$ obtained by BET was 119 m$^2$/g, which is comparable to that of a compact mesoporous TiO$_2$-nanocrystalline film (117 m$^2$/g), as shown in Figure S1. Figure 1e illustrates the engineering mechanism of absorbed light to maximize the intensity of the light. When compared to a conventional P25 electrode, wherein a large portion of light is not only reflected but also limited to having a single pass through the device, H-TiO$_2$ offers a much improved mechanism for light harvesting. Here, the double sided-large dimension structures function as a mirror at both ends: maximally absorbed photons by the help of the AR layer on the top are scattered further by the st-TIO frame on the bottom, permitting effective light harvesting via multi-passes of entrapped photons through the entire device.

The enhanced AR property and photon trapping effects due to coupling of properly designed double-sided light harvesting layers can be quantified by the (specular angle) reflection spectra and diffuse reflection spectra, respectively. The reflection spectra were obtained by shining light from a 5 µm-thick pure P25 and 5 µm top-surface patterned mesoporous P25 (p-P25) composed of 4.5 µm-thick P25 and an upper layer of patterns with a diameter of 500 nm and 500 nm thickness. P25 was chosen as a control sample since it has been shown to have excellent features as a working electrode in photovoltaic devices. In this configuration, the light passes through air/(TiO$_2$ patterns)/nanocrystalline TiO$_2$ interfaces from air to the TiO$_2$ nanocrystalline layer, and it is assumed that the majority of reflectance originates from the interface of air (n = 1) and TiO$_2$ (n = 2.49) due to the notable difference of the refractive index. Figure 2a shows that P25 gives ~20% lower reflectance as compared to the pure TiO$_2$ structure over the entire spectral range. This implies that much more light is able to participate in the conversion process in p-P25. This is attributable to the gradual change in the refractive index between the air and TiO$_2$ nanocrystalline layers in p-P25 (1/1.63/2.49, n$_{AR}$ = 1.63 as calculated in the Supporting Information), since the porosity of the patterned TiO$_2$ layer can modulate the effective refractive index of the diffraction grating layer. In order to investigate the intensity of scattered light under beam irradiation on the device, the diffuse reflectance spectra of 6 µm-thick TiO$_2$ films made in the form of P25, P25/st-TIO, and H-TiO$_2$ were measured as shown in Figure 2b. A schematic illustration of the samples is shown in Figure 2e. Interestingly, compared to the pristine P25 film, both films with a 1 µm-thick st-TIO layer (P25/st-TIO and H-TiO$_2$) showed significantly higher diffuse reflection spectra in the visible spectrum of light (380–800 nm), indicating that the absorbed light was efficiently scattered within the nanostructure film. The effect of introduction of the st-TIO layers into the P25 film was further exploited by the transmittance spectra by shining light from top to bottom (Figure 2c). As expected from the low absorption coefficient of P25 in the visible region, a significant fraction of the light passes through 6 µm-thick P25 at the longer wavelength range, which clearly shows light harvesting did not occur (black curve). However, when a 1 µm-thick st-TIO layer was introduced into 5 µm-thick P25 on the bottom, creating 6 µm-thick hybrid samples (blue and red curves), the transmission of light is greatly reduced, indicating that most of the light is re-absorbed by the structure in the devices. This is ascribed to the introduction of a few hundred nanometer scale structure as the photon trapping layer. Finally, UV-visible spectroscopic measurements were undertaken in order to explore the ultimate optical response of nanostructured TiO$_2$ with and without light harvesting layers. Figure 2d shows the UV-visible absorption spectra of an all 6 µm-thick plain P25, p-P25, P25/st-TIO, and H-TiO$_2$ illuminated at normal incidence from the air side. The absorption peaks spanning from 350 to 400 nm were assigned to the absorption of the conventional TiO$_2$ nanoparticles.
Figure 1 | Schematic illustrations and electron microscopy images of TiO₂ structures. (a) Schematic illustration of the fabrication process for the hybrid TiO₂ structure (H-TiO₂). (b) Top view SEM image of st-TiO for photon trapping. The inset shows transmittance electron microscopy (TEM) image of the mesoporous surface textured-TiO₂ (st-TiO₂). (c) Top view SEM image of the patterned mesoporous TiO₂ layer for anti-reflection. The inset shows the close-up image of the P25 particles. (d) Cross-sectional SEM image of H-TiO₂. (e) Schematic illustration of pathway of incoming light engineered by H-TiO₂ compared with the case of P25.
An intense and wider spectrum with a broad absorption band ranging from 400 to 800 nm compared to pure P25 is achieved when either a top or bottom light harvesting layer is adopted. Notably, st-TIO contributed more to increment of the absorption band than anti-reflective p-P25 layer, implying the greater influence of photon trapping effects than AR effects on light harvesting. Conclusively, when the double sided layer is applied, greatly enhanced UV spectra in the visible range are observed, even though the density of TiO2 particles is lower than that of P25. The enhanced UV absorption of H-TiO2 is attributed to the synergetic combination of AR and photon trapping effects obtained by the larger dimension nanostructures, which allow for longer effective light path lengths and thus create a clear and strong absorption band in the longer wavelength region ranging from 400 nm to 800 nm.

In order to further enhance the working capability of H-TiO2 in the visible wavelength region, CdSe quantum dots (QDs) were deposited onto the H-TiO2 structure. The deposition of CdSe on TiO2 was carried out via 20 cycles of successive ionic layer adsorption and reaction (SILAR) method. Figure 3a presents TEM images of the CdSe QDs deposited onto H-TiO2. Slightly randomly deposited spherical QDs with an average particle size of 4.5 nm with lattice spacings of 0.35 and 0.20 nm, which are consistent with the known values for the (002) and (103) planes of CdSe are observed. The enhanced optical properties due to the introduction of QDs on the hybrid samples are compared with the properties observed for CdSe/P25 and pure H-TiO2 in Figure 3b. An apparent increase and broadening in the optical density near the band gap of CdSe covering a range of 400–800 nm confirms the creation of CdSe nanoparticles and the ability of QD-sensitization. It should be noted that the overlap of the excitation wavelength of QDs and the regions of the wavelength with great absorption of H-TiO2 is crucial to maximize the efficient sensitization of QDs. CdSe/P25 also presents enhanced
absorption after the deposition of CdSe, yet with a much lower absorption coefficient than that of CdSe/H-TiO₂, likely due to low absorption of pristine P25 in the visible region where CdSe is excited. Figure 3c presents photoluminescence (PL) spectra of the CdSe QDs and CdSe/H-TiO₂. Compared to the PL spectrum of CdSe, which shows a typical near-band-edge emission at 630 nm, that of CdSe/H-TiO₂ is remarkably quenched. This clearly indicates a facile electron transfer mechanism at the interface between CdSe and TiO₂ nanoparticles, and thus low likelihood of electron-hole recombination within CdSe.

**Discussion**

In order to investigate the hydrogen generating efficiency of the samples, photo-electrochemical measurements were carried out on a P25 control sample, p-P25, P25/st-TiO₂, H-TiO₂, and CdSe/H-TiO₂ using Ag/AgCl, a Pt mesh, and as-prepared samples as a reference electrode, counter electrode, and working electrode, respectively, in a 0.24 M Na₂S and 0.35 M Na₂SO₃ electrolyte solution. This solution is commonly used as an electrolyte for quantum dot-based PEC systems, serving as the sacrificial agent. As shown in Figure 4a, which illustrates the electron transfer mechanism of CdSe/H-TiO₂, after the excited electrons are transferred from quantum dots to the conduction band of TiO₂, the electrons are withdrawn to the counter electrode where the hydrogen generation reaction occurs (An image capturing a large amount of H₂ gas generation is presented in Figure S6 along with a movie clip). Meanwhile, the holes diffused to the surface of TiO₂ or CdSe oxidize S²⁻. Figure 4b presents linear sweep voltammograms of the pristine P25, p-P25, P25/st-TiO₂, and H-TiO₂, which is described by the direct photoexcitation of electrons from the valence band of TiO₂ under AM 1.5 G simulated sunlight illumination (100 mW/cm²). All samples with a light harvesting layer show higher photocurrent density than the P25 control sample. Surprisingly, the photocurrent density of an optimized H-TiO₂ reaches ~1.65 mA/cm² at an applied bias of 0.5 V, which is approximately 2.5 times greater than that of the P25 control sample. This is the highest reported value of current density achieved using a single TiO₂ electrode material, implying that our approach to maximize light harvesting via engineering structural control is successful and also suggesting massive potential of pristine H-TiO₂ in photoconversion systems. Figure 4c compares linear sweep voltammograms of CdSe/H-TiO₂ and CdSe/P25. When the system is further supported by QD sensitization, the photocurrent density of CdSe/H-TiO₂ reaches ~16.2 mA/cm², which is 35% higher than that obtained with CdSe/P25, under illumination of 100 mW/cm². More surprisingly, as shown in Figure 4c, the photocurrent density of CdSe/H-TiO₂ created by only visible light illumination (blue curve, AM 1.5, a cutoff filter (>420 nm)) reaches ~14.2 mA/cm², ~88% of that of UV-visible irradiation (red curve) and comparable to the values under UV-visible irradiation in other detailed studies.
solely visible working capability of CdSe/H-TiO₂. This surprising result is ascribed to the large enhancement of absorption of H-TiO₂ in the visible range, which matches well with the excitation wavelength of CdSe, via efficient light harvesting. Figure 4d shows the photoresponse over time (J-t curve) of CdSe/H-TiO₂ measured at 0.5 V with chopped illumination of AM 1.5 at a rate of 10 s exposure followed by 10 s non-illumination for UV-visible and visible light.

The sharp spike in the photocurrent for both cases during the on/off illumination cycles indicates fast transport of photogenerated electrons from QDs to the current collector via the interior TiO₂ networks in accordance with the result of the PL study (Figure 3c). The greater values in the photocurrent density of CdSe/H-TiO₂ than in that of CdSe/P25 can be confirmed by a lower internal resistance measured by the electrochemical impedance spectroscopy (EIS). The
PEC measurements. The PEC performance of the TiO₂ electrodes was exploited in a three-cell electrode system under front-side illumination of AM 1.5 G (Newport solar simulator) using a potentiotstat (Princeton Applied Research VersaSTAT3, AMETEK). An Ag/AgCl electrode and a Pt mesh were used as reference and counter electrodes, respectively. A solution of 0.24 M Na₂S and 0.35 M NaSO₃ was used as an electrolyte and sacrificial solution. An exposed area of the working electrode was 0.25 cm². The photocurrent response was evaluated under chopped illumination from a Newport solar simulator equipped with an AM 1.5 G filter, calibrated with a standard Si solar cell simulating AM 1.5 G. The power of the solar simulator was measured to be 100 mW/cm². Photocurrent stability tests were carried out by measuring the photocurrent produced under chopped light irradiation (light/dark cycles of 10 s) at a bias of 0.5 V versus Ag/AgCl. The PEC performances in the visible range were acquired by the solar simulator coupled with a UV cutoff filter (λ > 420 nm) in Figure S5. EIS measurement were performed under the same condition of gas evolution under illumination of AM 1.5 G.

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Author contributions

The experiments were designed by K.K. and M.K. with suggestions from J.J.; K.K. prepared H-TiO2 samples. K.K. and M.K. performed PEC and optical property measurements. S.K. performed BET analysis and EIS measurements. J.J. analyzed the data and wrote the manuscript. All authors discussed the results and commented on the manuscript.

Additional information

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