Joule Heating and Viscous Dissipation on Effects on MHD Flow over a Stretching Porous Sheet Subjected to Power Law Heat Flux in Presence of Heat Source

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Abstract
In the present work we investigate the effects of Joule heating and viscous dissipation on MHD fluid flow. The viscous incompressible fluid flows over a stretching porous horizontal sheet subjected to power law heat flux in presence of heat source. The equations of momentum and heat transfer governing the problem are transformed into a system of dimensionless differential equations, which in turn solved numerically using shooting technique. The effects of the Joule heating parameter, permeability parameter, heat source parameter, Eckert number and Prandtl number are discussed and tabulated.

Keywords
Viscous Dissipation, Heat Transfer, Porous Medium, Heat Flux

1. Introduction
Studies of boundary layer flow of viscous fluid and heat transfer on a continuously moving surface are of great importance in view of their relevance to many technological processes, the manufacturing of aerodynamic extrusion of plastic sheets, cooling of metallic paths rapidly, many industries like fiber, glass and polymers. Sakiadis [1] was the first researcher who studied the boundary layer flow of a viscous fluid adjacent to a continuously stretching vertical sheet. A combined experimental and analytically study of the laminar and turbulent flow and temperature distribution in the boundary layer on a continuously moving surface has been carried out by Tsou et al. [2]. Erickson et al. [3] and Gupta et al. [4] extended Sakiadis problem, they taking into account suction or
blowing of the fluid at the moving surface. Since that many researcher extended the problem in various ways. Afzal et al. [5] and Banks [6] considered the power law stretching of the moving sheet. Magnetohydrodynamics (MHD) is the study of the magnetic properties of electrically conducting fluids, such as plasmas, liquid metals, and salt water or electrolytes. Hannes Alfven [7] was the first who studying MHD, He described the class of MHD waves now known as Alfven waves, for which he received the Nobel Prize in Physics in 1970. Recently, magneto hydrodynamic flow and heat transfer have been studied by many researchers such as Abo-Eldahab [8] and Chakarabati et al. [9]. Due to its ever increasing engineering applications such as the flow through packed beds, environmental pollution, blood theology and many other applications the extension of the problem to porous sheet attracted many researchers. Varjravelu et al. [10] studied the laminar flow and heat transfer of a viscous fluid over a stretching surface with viscous dissipation. Chamkha [11] studied the problem of thermal radiation effects on MHD forced convection laminar flow adjacent to a non-isothermal wedge in the presence of a heat source or sink. Abo-Eldahab et al. [12] studied the viscous dissipation and Joule heating effects on magneto hydrodynamic free convection from a vertical plate with power-law variation in surface temperature in the presence of Hall and ion-slip currents. Jaber [13] studied the Transient magneto hydrodynamic mixed double diffusive convection along a vertical plate embedded in a non-Darcy porous medium with suction or injection. Abo-Eldahab and Elbarbary [14] examined heat and mass transfer over a vertical plate in the presence of a uniform and strong magnetic field and Hall current. Abo-Eldahab and El Aziz [15] investigated the effects of Hall current and Joule heating on electrically conducting fluid past a semi-infinite plate with strong magnetic field and heat generation/absorption. Jaber [16] studied the effect of Hall currents, and variable viscosity on free convective flow past a semi-infinite continuously stretching plate in presence of radiation. Jaber [17] studied the effect of Hall currents and variable fluid properties on magneto hydrodynamics flow past a continuously stretching vertical plate in the presence of radiation. Viscous dissipation effect plays an important role in natural convection in various devices which are subjected to large variations of gravitational force or which operate at high speeds Jaber [18] examined the effects of viscous dissipation and joule heating on magneto hydrodynamics flow of a fluid with variable properties past a stretching vertical plate. Mandal, et al. [19] studied the effect of Hall current on MHD couette flow between thick arbitrarily conducting Plates in a rotating system. Heat flux or thermal flux is the rate of heat energy transfer through a given surface per unit time Mandal et al. [20] study the boundary layer flow and heat transfer towards an exponentially stretching porous sheet embedded in a porous medium with variable surface heat flux.

The purpose of this present work is to study the combined effects of Joule heating and viscous dissipation on laminar flow and heat transfer of a viscous fluid over a horizontal porous sheet subjected to power law heat flux in presence of heat source, also the plate is continuously stretching. The governing equations are transformed into dimensionless ordinary differential equations. Numerical calculations up to desired level
of accuracy were carried out to investigate the effects of all parameters involved in the problem under consideration to illustrate the results graphically.

2. Mathematical Formulation

The present work generalizes the problem of heat transfer over a stretching porous sheet subjected to power law heat flux in presence of heat source which presented by Hitesh Kumar [21]. Consider the laminar steady two dimensional flow of an incompressible viscous fluid over a horizontal plate coinciding with the plane \( y = 0 \). A uniform strong magnetic field of strength \( B_0 \) is subjected normal to the plate. The plate is subjected to a power law heat flux in presence of heat source. The plate is continuously stretching by two equal powers with opposite directions. The speed is assumed to be proportional to its distance from the slit point. The flow is considered to be in the \( x \)-direction which is taken along the plate, see Figure 1.

Under these assumptions and taking into account the Boussinesq approximation, the boundary layer laminar flow is governed by the following equations:

\[
\begin{align*}
\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} &= 0 \quad (1) \\
\frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} &= \frac{\nu}{\rho} \left( \frac{\partial^2 u}{\partial y^2} - \frac{\sigma B_0^2}{\rho} u \right) \quad (2) \\
\frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} &= \frac{k}{\rho C_p} \frac{\partial^2 T}{\partial y^2} - \nu \left( \frac{\partial u}{\partial y} \right)^2 + \frac{\sigma B_0^2}{\rho C_p} u^2 \quad (3)
\end{align*}
\]

The physical problem suggested the following initial and boundary conditions:

\[
\begin{align*}
\text{at } y = 0 : u &= cx, v = -v_0, \frac{\partial T}{\partial y} = Bx^2 \\
\text{at } y \to \infty : u &= -cx, T = T_\infty
\end{align*}
\]

where \( c \) is the stretching rate which considered to be a positive integer.

The continuity equation suggested the solution

\[
\begin{align*}
u &= cx f'(\eta), \quad v = -\sqrt{\nu c f} \left( \frac{\eta}{\nu c} \right), \quad \eta = \sqrt{cu} y, \quad T = T_\infty + Bx^2 \sqrt{\nu c} \theta(\eta)
\end{align*}
\]

Using these dimensionless parameters the governing equations, Equations (2) and (3) can be reduced to the following system of ordinary differential equations

\[
\begin{align*}
 f'''' + ff'' - f'^2 - Mf' &= 0 \quad (5) \\
 \theta'' + Pr f' \theta - 2 Pr f' \theta - \beta \theta + Pr Ec f'^2 + Ec f'^2 &= 0 \quad (6)
\end{align*}
\]

![Figure 1. Physical coordinate system.](image-url)
The boundary conditions are transformed to:

At the surface of the plate \( f'(0) = 1, f(0) = \frac{v_0}{\sqrt{\eta}}, \theta'(0) = 1 \)

Also, as \( \eta \to \infty : f' = f = \theta = 0 \)

where \( Pr = \frac{\rho \nu C_p}{k} \) is the Prandtl number, \( Ec = \frac{C_p}{\nu C_s \sqrt{\nu}} \) is the Eckert number

\( J = \frac{\sigma B^2 \nu C_p}{k} \) is the Joule heating parameter, \( \beta = \frac{\nu Q}{ck} \) is the heat source parameter.

The physical quantities of interest Nusselt number and the local shear stress \( \tau_w \) are defined as

\[ q_w = -k \left( \frac{\partial T}{\partial y} \right)_{y=0} \]

Thus the Nusselt number is given by

\[ Nu = \frac{q_w x}{k} = -Bx^2 \theta'(0) \]

The wall shear stress is defined as

\[ \tau_w = \mu \left( \frac{\partial u}{\partial y} \right)_{y=0} = \frac{\mu x c^2}{\sqrt{\nu}} f'(0) \]

The numerical values of \( f'(0) \) and \( \theta'(0) \) for several values of the Magnetic number, Joule heating parameter, heat source parameter, Eckert number and Prandtl number are tabulated in Table 1.

**Table 1.** Variation of the dimensionless skin friction and local Nusselt number for various parameters.

| \( Pr \) | \( Ec \) | \( M \) | \( J \) | \( \beta \) | \( -f(0) \) | \( \theta(0) \) |
|---|---|---|---|---|---|---|
| 0.72 | 3 | 1 | 1 | 1 | 2.00007 | 0.719796 |
| 2 | 3 | 1 | 1 | 1 | 2.00007 | 1.26137 |
| 5 | 3 | 1 | 1 | 1 | 2.00007 | 1.58979 |
| 0.72 | 5 | 1 | 1 | 1 | 2.00007 | 1.55876 |
| 0.72 | 8 | 1 | 1 | 1 | 2.00007 | 2.8172 |
| 0.72 | 3 | 3 | 1 | 1 | 2.56155 | 1.00835 |
| 0.72 | 3 | 5 | 1 | 1 | 3 | 1.25537 |
| 0.72 | 3 | 1 | 3 | 1 | 2.00007 | 1.36839 |
| 0.72 | 3 | 1 | 5 | 1 | 2.00007 | 2.01699 |
| 0.72 | 3 | 1 | 1 | 3 | 2.00007 | 0.437762 |
| 0.72 | 3 | 1 | 1 | 5 | 2.00007 | 0.315727 |
Numerical Technique

The transformed governing system consists of two non-linear ordinary differential equations. They were solved numerically up to desired level of accuracy using the fourth order Runge-Kutta method with the shooting technique. The effect of all the involved parameters on the velocity and temperature profiles, the local Nusselt number and the shear stress is presented and discussed. The effects of the involved parameters were presented graphically as shown in Figures 2-7. Figure 2 and Figure 3 present the effect of the magnetic parameter $M$ on the velocity and the temperature profiles, the increasing of magnetic parameter $M$ causes decreasing in the velocity and give rise to increase the temperature profiles. Figure 4 shows that the increasing of the Eckert number leads to increase the temperature profiles. The effect of the Prandtl number is presented in Figure 5, it is obvious from the figure that the increasing in the Prandtl number leads to a markedly increasing in the initial magnitude of the temperature in the boundary

![Figure 2](image2.png)

**Figure 2.** Variation of dimensionless velocity $f'$ for various values of magnetic number $M$ with $Pr = 0.72, \beta = 1, Ec = 3, J = 3$.

![Figure 3](image3.png)

**Figure 3.** Variation of dimensionless temperature $\theta$ for various values of magnetic number $M$ with $Pr = 0.72, \beta = 1, Ec = 3, J = 3$. 
Figure 4. Variation of dimensionless temperature $\theta$ for various values of Eckert number $Ec$ with $Pr = 0.72$, $\beta = 1$, $J = 3$, $M = 1$.

Figure 5. Variation of dimensionless temperature $\theta$ for various values of $Pr$ and Pr number $Pr$ with $\beta = M = 1$, $J = 3$, $Ec = 3$.

layer closed to the plate while this effect takes the reverse order as the fluid goes far from the plate i.e. temperature profile decreases as the fluid get far away from the plate. Figure 6 investigates the effect of the Joule heating parameter, the increase in the Joule heating parameter causes markedly increase in the magnitude of the temperature in the boundary layer. Finally, Figure 7 presents the effect of the heat source parameter $\beta$ on temperature profile, it is clear that the increasing in the heat source parameter $\beta$ leads to decrease the temperature in the boundary layer. Table 1 shows that the local skin friction for the flow decreases as the magnetic parameter $M$ increases, while other parameters do not affect the local skin friction. Also, it is easily seen that the increasing in the values of $Pr$, $Ec$, $M$ and Joule heating parameters increase the local Nusselt number, while it is decreased by the increasing of the heat source parameter $\beta$. 

Table 1
Figure 6. Variation of dimensionless temperature $\theta$ for various values of Joule heating parameter $J$ with $Ec = 3$, $Pr = 0.72$, $\beta = M = 1$.

Figure 7. Variation of dimensionless temperature $\theta$ for various values of heat source parameter $\beta$ with $Ec = J = 3$, $Pr = 0.72$, $M = 1$.

3. Conclusions and Remarks

The combined effects of viscous dissipation and Joule heating on the steady, laminar, convection heat transfer over a horizontal plate is investigated, the plate is permeable and continuously stretching with a power-law heat source. The governing equations are transformed into a system of ordinary differential equations that solved numerically. Results are tabulated and presented graphically for various values of the involved parameters.

The present study shows that:

1) Increasing the magnetic parameter $M$ increases the temperature and the local Nusselt number and decreasing the velocity and the local skin friction.

2) Increasing the Eckert number increases both the local Nusselt number and tem-
perature profiles.

3) Increasing the Prandtl number $Pr$ increases both the initial values of temperature and the local Nusselt number.

4) As expected the increasing in the Joule heating parameter increases the temperature profile and local Nusselt number, also it does not affect the skin friction.

5) As the heat source parameter $\beta$ increases, the temperature and local skin friction decrease.

6) No effect of $J$, $Ec$, $Pr$ and heat source parameter $\beta$ on velocity profiles so no figures are presented herein to limit the figures number, also they do not affect the local skin friction.

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Nomenclature

| Symbol | Definition |
|--------|------------|
| $u$    | Velocity component in x-direction |
| $v$    | Velocity component in y-direction |
| $x$    | Stream wise coordinate |
| $y$    | Direction normal to the plate |
| $Ec$   | Eckert number |
| $Pr$   | Prandtl number |
| $f$    | Dimensionless stream function |
| $g_p$  | Specific heat at constant pressure |
| $k$    | Thermal conductivity |
| $Nu$   | Nusselt number |
| $\tau$ | Shearing stress |
| $T$    | Temperature |
| $q_w$  | The local heat flux |
| $c$    | The stretching rate |
| $K$    | Permeability of the porous medium |
| $J$    | Joule heating parameter |
| $M$    | Magnetic number |
| $\xi$  | Differentiation with respect to $\eta$ only |

Greek Symbols

| Symbol | Definition |
|--------|------------|
| $\beta$ | Heat source parameter |
| $\eta$ | Pseudo similar variable |
| $\mu$  | Dynamical viscosity |
| $\rho$ | Density |
| $\nu$  | Kinematical viscosity |
| $\theta$ | Dimensionless temperature |

Subscripts

| Symbol | Definition |
|--------|------------|
| $w$    | Property at the wall |
| $\infty$ | Free stream condition |

Superscripts

| Symbol | Definition |
|--------|------------|
| $M$    | Magnetic number |

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