The information-theoretic view of quantum mechanics and the measurement problem(s)

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Abstract
Until recently Jeffrey Bub and Itamar Pitowsky, in the framework of an information-theoretic view of quantum mechanics, claimed first that to the measurement problem in its ordinary formulation there correspond in effect two measurement problems (simply called the big and the small measurement problems), with a different degree of relevance and, second, that the analysis of a quantum measurement is a problem only if other assumptions – taken by Pitowsky and Bub to be unnecessary ‘dogmas’ – are assumed. Here I critically discuss this unconventional stance on the measurement problem and argue that the Bub-Pitowsky arguments are inconclusive, mainly because they rely on an unwarranted extension to the quantum realm of a distinction concerning the foundations of special relativity which is in itself rather controversial.

Keywords Information-Theoretic View of Quantum Mechanics · Measurement Problem · Principle vs. Constructive Theories · Dynamics vs. Kinematics · Special Relativity

1 Introduction

In retrospect, it is far from surprising that in one of the great scientific works of the XXth century, the 1932 von Neumann book on the mathematical foundations of quantum mechanics (QM), an entire chapter is devoted to the problem of how to construct an ideal quantum–mechanical model of a measurement (von Neumann 1955, chapter VI). The von Neumann treatment, and the place occupied by this problem in his first formally rigorous formulation of quantum theory, already revealed how controversial the status of measurement in quantum mechanics would have been, to the extent that the very notion of measurement would turn out to be the locus classicus

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for emphasizing the lack of consensus on the interpretation of the theory. In fact, that chapter happens to be the major source of what is usually defined as the measurement problem: unlike the case of the vaguely defined classical-quantum distinction previously advocated by Bohr, von Neumann explicitly confronts the implications of the assumption that – in the context of a measurement of a physical quantity on a quantum system $S$ with an apparatus $A$ – the laws of QM govern both $S$ and $A$. As a matter of fact, the measurement problem is widely taken to be a true touchstone for a classification of the main different interpretations of QM. If, according to the folklore, a standard measurement procedure induces a ‘collapse’ of a superposition state of the joint system $A + S$ into one of its components, the disagreement arises at the starting block: is this ‘collapse’ a real physical process or not? The most radical No-Answer leads to the claim that such ‘collapse’ is just a sort of perspectival effect (Everett-style of thinking). On the other hand, the definite Yes-Answer leads to the assumption of new laws to be added to standard QM, laws that dictate the how and when of a physical collapse, so as to make it compatible with the empirical fact that at the end of a measurement we obtain a definite outcome (GRW-style of thinking). A sort of Yes-and-No-Answer is provided by Bohmian mechanics: in this theory, given the joint system $A + S$ as a subsystem of the entire universe $U$, it is possible to naturally associate with $S$ a so-called conditional wave function; both this wave function and the wave function of the $A + S$ joint system are subject to genuine collapses in accordance with the Born probability rule – as far as they are both subsystems of $U$ – whereas the wave function of $U$ is not.1

Independently of the problem of how the collapse of the post-measurement state of $A + S$ is to be interpreted, though, there are those who take a different approach and question the very idea of a real measurement problem in its ordinary formulation. In this approach, developed in more recent times especially by the late Pitowsky and Jeffrey Bub on the background of the so called ‘information-theoretic’ view of QM, it is claimed first that to the measurement problem in its ordinary formulation there correspond in effect two measurement problems (simply called the big and the small measurement problems), with a different degree of relevance and, second, that the analysis of a quantum measurement is a problem only if other assumptions – taken by Pitowsky and Bub to be unnecessary ‘dogmas’ – are assumed.

In the present paper I will provide a critical assessment of this unconventional stance on the measurement problem, and I will argue that the Bub-Pitowsky position is unsatisfactory, both logically and conceptually. First, in Section 2, I will question the alleged logical dependence of the measurement problem from two assumptions

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1 “This is the sense in which the pilot-wave theory explains the collapse postulate, and in general the rules that are instead postulates for ordinary QM” (Norsen, 2016, p. 6). As is well known, the number of interpretation largely exceeds the list including Everett, Bohm and GRW, but I take these to be the most transparent in taking an unambiguous stance on the issue of the exact nature of collapse. Moreover, also Everettians and Bohmians can be said to ‘add’ something to the QM description, although in a more nuanced way, given that their view are predictively equivalent to standard QM: the Everettians add an explanation of why we perceive just one world, whereas the Bohmians add information on the position of the quantum systems under scrutiny, information that is supposed to be unavailable in the standard formulation.
(considered by Bub-Pitowsky as ‘dogmas’), arguing that the measurement problem in fact depends on neither (that is, I can question both assumptions and still I have to face the measurement problem). I will investigate the very definition of the big and the small measurement problems provided by Pitowsky and Bub, in order to clarify their logical status and mutual relationships. Since Pitowsky and Bub take the status of these problems \textit{qua} problems to depend on two claims that, in turn, they take to be true dogmas of the folklore view of QM, an integral part of my analysis will be to clarify how the big/small distinction fares with respect to these two ‘dogmas’. I will subsequently focus on the \textit{conceptual} status of the Bub-Pitowsky proposal for a (dis) solution of the measurement problem. In Section 3, I will first review how this proposal is based on the extension to the quantum realm of a specific interpretation of the Minkowskian geometry of special-relativistic spacetime, whereas in Section 4 I will argue that that interpretation is controversial: as a consequence, its extension to quantum mechanics can also be questioned, together with proposal the very plausibility of the whole Bub-Pitowsky strategy of (dis)solution of the measurement problem. It must be said that the Bub-Pitowsky is not the last word on the issue within the information-theoretic view of QM – and in the last part of the section I will deal with a successive formulation, developed in particular in Bub, 2016 and Bub, 2018. Still, I take the Bub and Pitowsky, 2010 to be a crisp and well-characterized defense of an information-theoretic view on the measurement problem, so as to deserve an investigation on its own; moreover, the arguments that I put forward in my critical evaluation of the Bub-Pitowsky proposal are markedly different from those developed in the existing literature on it. In Section 5 I will deal with a further aspect of the Bub-Pitowsky (dis)solution proposal, connected with the extent to which alternative interpretations of QM should be accepted even if they do not provide new empirical predictions. My overall conclusions will be drawn in the final Section 6.

2 How many measurement problems are there? A logical analysis

Jeffrey Bub and Itamar Pitowsky initially present the measurement problem in a (rather) standard way:

The measurement problem is the problem of explaining the apparently ‘irreducible and uncontrollable disturbance’ in a quantum measurement process, the ‘collapse’ of the wavefunction described by von Neumann’s projection postulate.” (Bub & Pitowsky, 2010, p. 438).

As we will see later, Bub and Pitowsky support an information-theoretic approach to QM but with a realistic tone, and as a consequence they are unhappy with a purely instrumentalistic reading of the collapse. In this vein, they claim that the ‘irreducible and uncontrollable measurement disturbance’ – vaguely associated with collapse in a Copenhagenish style – fails to receive a decent explanation:

Without a dynamical explanation of this measurement disturbance, or an analysis of what is involved in a quantum measurement process that addresses the issue […], the theory qualifies as an algorithm for predicting the probabil-
ties of measurement outcomes, but cannot be regarded as providing a realist account, in principle, of how events come about in a measurement process. (Bub & Pitowsky, 2010, p. 435).

This statement might be taken simply as a rather common manifestation of dissatisfaction with instrumentalism in the interpretational debate on QM, were it not for the dynamical qualification of the explanation that Bub and Pitowsky take to be necessary. This reference, to which we will come back later, is also explicit in the claim that there is no just one measurement problem, but in fact two:

The ‘big’ measurement problem is the problem of explaining how measurements can have definite outcomes, given the unitary dynamics of the theory: it is the problem of explaining how individual measurement outcomes come about dynamically. The ‘small’ measurement problem is the problem of accounting for our familiar experience of a classical or Boolean macroworld, given the non-Boolean character of the underlying quantum event space: it is the problem of explaining the dynamical emergence of an effectively classical probability space of macroscopic measurement outcomes in a quantum measurement process. (Bub & Pitowsky, 2010, p. 438, my emphasis)

As is clear from the semantics of the ‘big/small’ distinction, the small measurement problem is taken to be relatively easy to solve. Putting aside for a moment the issue of characterizing different layers of the natural world as ‘Boolean’ or ‘non-Boolean’ (a far-from-innocent issue, to which we will return later in connection with the ‘big’ measurement problem), Bub and Pitowsky do not diverge from the mainstream approach where decoherence does the job:

The ‘small’ measurement problem is resolved by considering the dynamics of the measurement process and the role of decoherence in the emergence of an effectively classical probability space of macroevents to which the Born probabilities refer.” (Bub & Pitowsky, 2010, p. 438).3

The heart of the matter lies in the ‘big’ measurement problem, whose status is controversial according to Bub and Pitowsky:

The big measurement problem depends for its legitimacy on the acceptance of two dogmas. […] The first dogma is Bell’s assertion that measurement should never be introduced as a primitive process in a fundamental mechanical theory like classical or quantum mechanics, but should always be open to a complete

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2 Brukner 2017 also proposes a view according to which the measurement problem, as ordinarily understood, in fact splits into two versions. The Brukner overall approach to the solution to the measurement(s) problem, however, has a ‘relational’ flavour and significantly diverges from the Bub-Pitowsky approach. More recently, Alexander Meehan proposed a new measurement-related foundational problem, that he calls the control problem (Meehan 2019): the proposal has given rise to criticisms by Vaidman 2020, to which Meehan 2021 replied. An overall assessment of the debate is provided by Hetzroni 2021.

3 It is widely believed that also this solution is a for-all-practical-purposes solution, at least in the standard presentation in which there is just one measurement problem (for a recent review Bacciagaluppi 2020).
analysis, in principle, of how the individual outcomes come about dynamically. The second dogma is the view that the quantum state has an ontological significance analogous to the ontological significance of the classical state as the ‘truthmaker’ for propositions about the occurrence and non-occurrence of events, i.e., that the quantum state is a representation of physical reality.” (Bub & Pitowsky, 2010, p. 438, my emphasis).

In the present section, I will take for granted the Bub-Pitowsky formulation of the big measurement problem and I will consider its logical relation with the two ‘dogmas’, independently from the details of that formulation, whereas I will focus on the controversial status of the formulation by the conceptual point of view in the next two sections.

A preliminary remark concerns the status of ‘dogma’ that Bub and Pitowsky ascribe to the two assumptions that, in their opinion, would be the ground for the emergence of the measurement problem. According to the Oxford Dictionary, a dogma is “a belief or set of beliefs held by a group or organization that others are expected to accept without argument”: if the ‘group’ in question is the scientific community involved at different degrees with the foundations of quantum mechanics, none of the two assumptions emphasized by Bub and Pitowsky can be easily recognized as a dogma in this sense. As to the first – concerning the primitive or derivative status of the notion of quantum measurement – the Bell meta-theoretical stance is simply a legitimate assumption that, like many others on the market of ideas, comes with its possible strengths and weaknesses, and it is simply false that the scientific community in the foundations of quantum mechanics ‘accepts it without argument’. The very fact that the problem of coexistence of unitary and non-unitary dynamics, recalled above, has been widely acknowledged as a crucial issue – if not the issue – in the foundations of quantum mechanics should suggest that the Bell view is indeed a serious option, but to turn it into a ‘dogma’ is no argument, on the contrary it is an exercise in dogmatism itself. Something similar can be said about the second assumption, concerning the true ontological character of the notion of state: what the notion of quantum state stands for exactly is matter of debate since the early days of quantum mechanics, and its representation in realistic terms – whatever realistic might mean – can hardly be taken as a widely and acritically shared view in the quantum foundations’ community. Therefore, in the rest of the section, I will leave aside the original ‘dogmatic’ qualification given by Bub and Pitowsky, referring instead to the first ‘dogma’ with the expression <Measurement-Derivative (MD) assumption>, and to the second ‘dogma’ with the expression <States-As-Entities (SAE) assumption>.

Let us come now to the relation between these two assumptions and the measurement problem. The alleged dependence of the big measurement problem on the (MD) assumption is unconvincing for the very basic reason that the rejection of MD and the resulting inclusion of measurement among the primitive theoretical notions do not explain away the big measurement problem per se. The vaguely defined coexistence of unitary and non-unitary dynamics has been considered puzzling and unsatisfactory since the origins of QM, quite independently from whether measurement should have a primitive or derivative status in a fundamental theory such as
QM. Therefore, even if we decide to drop the MD assumption, this does not put us in any better position to solve, or dissolve, the big measurement problem.

The dependence of the big measurement problem on the SAE assumption seems even more problematic. In an earlier paper, Pitowsky had described the assumption as a requirement according to which:

the quantum state is a real physical entity, and that denying its reality turns quantum theory into a mere instrument for predictions. *This last assumption runs very quickly into the measurement problem.*” (Pitowsky, 2006, p. 214, my emphasis).

Since, according to Pitowsky, “the BIG problem concerns those who believe that the quantum state is a real physical state which obeys Schrödinger’s equation in all circumstances” (Pitowsky, 2006, p. 232, capital in the original text), Pitowsky’s point seems to be that the big measurement problem is a problem really only if we assume quantum states as real entities. This point makes the ‘dogmatic’ qualification sound less unreasonable than the case with the first ‘dogma’, but on what basis can we argue that it is the assumption of quantum states as ‘real entities’ that leads us to require from quantum mechanics a dynamical description of the measurement?

And what should we exactly mean when we say that a state is ‘a real entity’?

Let me first focus on the very problem of interpreting a physical state as a ‘real-thing-out-there’. With the aid of a classical-sounding language Bub and Pitowsky depict the second ‘dogma’ as the extension to the quantum realm of an assumption taken to be obviously unproblematic in the realm of classical physical theories, namely that in these theories a physical state is a ‘real-thing-out-there’. Moreover, they seem to assume that, for a state of a classical theory, to be a “truthmaker for propositions about the occurrence and non-occurrence of events” or a “representation” of physical reality is equivalent to be a ‘real-thing-out-there’. As a matter of fact, things do not seem so straight: even if we put aside the remark that one thing is to say that a state is a real entity and quite another to say that ‘it represents something in physical reality’, also in classical theories the relation between a ‘state’ (according to a given theoretical framework) and the domain of ‘real-things-out-there’ – be they medium-size objects or macroscopic events or properties – is complex and far from direct.

Let us consider briefly a classical-mechanical framework. In this case according to the usual intuition, well-entrenched into the formal detailed development of any such framework, the objects the theory is about can be considered as real entities, endowed with well-defined physical properties that can be easily imagined as possessed properties, quite independently from any attempts on our part to check the possession of such properties on an experimental basis. Already in this framework, states can be conceived as descriptions of the ways in which things-out-there stand. Namely, at a time \( t \) we assume a classical physical system \( S \) (for simplicity, a Newtonian one-dimensional point-particle) to be in a given, conventional state represented by a pair of values for position \( x \) and momentum \( p \) and all remaining significant properties of \( S \) depend on the values of \( x \) and \( p \). The set of all such points, endowed with a suitable geometric structure, is the phase space: all the physical quantities that are assumed to be relevant to \( S \) (the classical observables for \( S \)) are introduced.
as continuous, real-valued functions on its phase space and the theory provides the formal recipe for describing the dynamics of the system. If, for whatever reasons, we are unable to specify exact values for $x$ and $p$ at a certain time, we describe the state of the system via a probability density $\rho(x,p)$, for which too – via the Liouville equation – a dynamics is secured. Therefore, if these are the ordinary intuition and the formal implementation of the notion of state in a classical mechanical framework, we may ask ourselves: did we somewhere need to assume that a state is – or needs to be – an entity? What kind of feature is such entity-language supposed to pick out exactly? The above remarks on the role of the notion of state in a classical setting, a setting in which the intuition of pre-existing entities whose properties are independent from our attempts to have access to them is unproblematic, suggest that we need not require from states to be ‘entities’ in their own right.\footnote{Even less so in a classical statistical mechanical framework, in which a key role is played by the distinction between the macro-state and micro-state of a system, in view of the explanation of the irreversible thermodynamic behavior of macroscopic systems.} We do not at the intuitive level – in which states need not be conceived as physical entities in themselves, but rather descriptions of ways of being of physical entities\footnote{See Sudbery 2017 for a sustained defense of the claim that a state is not a thing but rather a description of things, in the context of temporal logics in a quantum framework.} – and we do not at the mathematical level – in which states are abstract, mathematical objects whose status of ‘entities’ is at best controversial and, in any case, in need of an engaging, Platonic-sounding argument in support.\footnote{I might say that I support a sort of ‘anti-Tractarian’ view of states (both classical and quantum). If, starting from the 1.1 Satz of the Tractatus, Wittgenstein appears to attribute reality to actual states of affairs (the ‘facts’), denying it to concrete particulars, I support the opposite view: we start from individual physical systems, whose possession of properties is described via the notion of state.}

Therefore, if the very assumption of states as real entities, in addition to its lack of clarity, appears to be unnecessary already in a non-quantum framework, it turns out to be even more dubious that such assumption is in fact assumed, explicitly or implicitly, in the quantum realm. On the contrary, the Pitowsky-Bub viewpoint not only takes this assumption to be a ‘dogma’, namely a statement that is endorsed uncritically, but also claims that it is only such an assumption that generates the ‘big’ measurement problem. Let us see, then, why the standard formulation of the ‘big’ measurement problem nowhere requires any assumption on the entity status of quantum states and why, as a consequence, the emergence of the ‘big’ measurement problem can be safely taken to be independent from the issue of the ‘reality’ of states.

In that formulation,\footnote{I refer here to the neat presentation in Maudlin 1995.} we assume quantum mechanics to describe measurement as a special kind of interaction, such that with the coupling $<\text{measured system} + \text{measuring apparatus}>$ determines a joint system whose states are supposed to evolve according to the main dynamical law of the theory, i.e. the Schrödinger equation (at least, up to the time of the measurement). Since in a measurement we are supposed to record an outcome for a physical quantity (which is well-defined for the measured system at hand), there are two possible scenarios:
(i) if the measured system’s state is an eigenstate of the physical quantity to be measured, the state of the joint system will be a state in which the component referring to the measuring apparatus will be unequivocally associated to the reading of (eigen)value of the quantity pertaining the measured system;

(ii) if the measured system’s state is not an eigenstate of the physical quantity to be measured, the state of the joint system will be a superposition, each component of which will be the product of measured system’s state and the measuring apparatus’ state, each corresponding to one of the possible (eigen)values for the physical quantity.

Now in the (ii) situation, namely when the measured system is in a superposition before the measurement takes place, the measurement problem amounts exactly to the fact that the following conditions cannot hold together:

C – The wave-function associated to the state of a system is a complete description of the state itself, namely there can be no finer specification of the properties that the system can exhibit in the event of a measurement;

L – The wave-function associated to the state of a system always evolves according to the Schrödinger equation;

D – Measurements always provide have determinate outcomes, namely at the end of the measurement process the measuring apparatus is found to be in a state that indicates which among the possible values turns out to be the outcome of the process itself.

Even if we suppose, for the sake of the argument, to distinguish between a state and a wave function, taken as the mathematical object that ‘describes’ the state, still the above argument nowhere depends on the assumption that either the wave functions or the states are real entities.⁸

³ A ‘kinematical’ solution to the measurement problem?
A conceptual analysis

I have argued above that the dependence of the big measurement problem on the two dogmas is controversial at best, but now let me focus on the very formulation of such problem by Bub and Pitowsky. When they say that.

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⁸ A caveat should be made here. The above discussion on states as ‘real’ entities has been conducted with reference to a rather intuitive view of what it means for something to be ‘real’. Namely, no specification of the issue has been attempted with the use of the conceptual resources of the so-called ontic vs. epistemic models of quantum states. For a recent paper that assesses the Pitowsky view on the issue on the background of those models, see Ben-Menahem 2020, whereas a recent evaluation of the Harrigan & Spekkens 2010 categorization and its relation with the map of the different interpretation of quantum mechanics is provided in Oldofredi & Lopez 2020. Personally, I tend to think that the notion of ontic state might be subject to criticisms analogous to that mentioned above, but this is a topic for a different paper.
the ‘big’ measurement problem is the problem of explaining how measurements can have definite outcomes, given the unitary dynamics of the theory: it is the problem of explaining how individual measurement outcomes come about dynamically (Bub & Pitowsky, 2010, p. 438)

they refer to a peculiar interpretation of ‘dynamics’. The conceptual framework for this interpretation, apparently independent from the description of any dynamical process per se governing the evolution of quantum states over time, relies on two major elements:

(i) an information-theoretic view of quantum mechanics, in which the latter should be “viewed not as first and foremost a mechanical theory of waves and particles […] but as a theory about the possibilities and impossibilities of information transfer.” (Clifton, Bub, & Halvorson, 2003, p. 1563, CBH, 2003 from now on);

(ii) on the background of (i), the adoption and extension to the quantum realm of the distinction between ‘kinematical’ and ‘dynamical’, as introduced in the debate on the foundations of special relativity.

I will first outline the Bub-Pitowsky information-theoretic (IT-) view of quantum mechanics in a very sketchy way that is useful to our purpose, and after I will pass on to the analysis of the extended formulation of the pair ‘kinematical’/ ‘dynamical’ that Bub-Pitowsky take to be relevant for their IT-view of quantum mechanics. This extended formulation is claimed to contribute to the (dis)solution – within the IT-view of QM – of what Bub-Pitowsky call the big measurement problem: I will argue in the Sect. 4 that such formulation is based on a controversial interpretation of of the pair ‘kinematical’/ ‘dynamical’ itself in special relativity and that, as a consequence, the ‘big’ measurement problem fails to be solved simply by the adoption of a Hilbert space structure with its ‘kinematical’ constraints. In fact, this interpretation assumes a certain kind of explanatory priority of the ‘kinematical’ over the ‘dynamical’ that is unwarranted: if we have reason to argue that this assumption fails to be convincing, then also the explanatory power of the non-Booleanity of the Hilbert space as-a-kinematical-framework can be questioned.

What is nowadays known as the information-theoretic view of quantum mechanics is not a monolithic view, and a detailed analysis of its possible, and up-to-date, variants is out of the scope of the present paper.9 The IT-view of QM certainly provided a new twist to what Henderson called the ‘reaxiomatisation programme in quantum mechanics’ namely “a programme which aims to reaxiomatise the theory in terms of postulates which are clearer, more ‘reasonable’ and more physically motivated.” (Henderson, 2020, p. 292). In this spirit, the IT-view of QM is in general lines an attempt to ground the interpretation of quantum mechanics on principles of an informational character, principles that are taken to be “fundamental information-theoretic ‘laws of nature’ “ (CBH, 2003, p. 1562), although the extent to which these principles should

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9 The most recent book-length philosophical defense of the IT-view of quantum mechanics is Janas et al., 2022.
be interpreted in realistic terms, i.e. as principles that informationally constrain the very nature of physical reality, is a matter of dispute (see Dunlap, 2022 on the point). The contribution of the CBH work to the development of the IT-view was the possibility to show that, on the basis of these principles, standard Hilbert space QM could be truly derived:

What CBH showed was that one can derive the basic kinematic features of a quantum-theoretic description of physical systems in the above sense from three fundamental information-theoretic constraints: (i) the impossibility of superluminal information transfer between two physical systems by performing measurements on one of them, (ii) the impossibility of perfectly broadcasting the information contained in an unknown physical state (for pure states, this amounts to “no cloning”), and (iii) the impossibility of communicating information so as to implement a certain primitive cryptographic protocol, called “bit commitment,” with unconditional security. They also partly demonstrated the converse derivation, leaving open a question concerning nonlocality and bit commitment. This remaining issue has been resolved by Hans Halvorson, so we have a characterization theorem for quantum theory in terms of the three information-theoretic constraints. (Bub, 2005, pp. 549-550)

To be true, the program to provide the theory with an axiomatic structure based on a restricted set of “more ‘reasonable’ and more physically motivated postulates” is far from new. The so-called quantum-logical approach to QM was developed already in the Sixties of the XXth century exactly with that aim, on the basis of the circumstance (first recognized in von Neumann 1932 and then developed in Birkhoff & von Neumann, 1936) that standard Hilbert space QM provides for the set of ‘yes–no experiments’ performable on a quantum system a non-classical algebraic structure – the structure $\mathbb{L}(H)$ of orthomodular (non distributive) lattice of projection operators acting on the Hilbert space $H$ associated to the quantum system in question. As a recent review on the issue clarifies,

there […] remains the question of why the logic of measurement outcomes should have the very special form $\mathbb{L}(H)$, and never anything more general. This question entertains the idea that the formal structure of quantum mechanics may be uniquely determined by a small number of reasonable assumptions, together perhaps with certain manifest regularities in the observed phenomena. This possibility is already contemplated in von Neumann’s Grundlagen (and also his later work in continuous geometry), but first becomes explicit—and programmatic—in the work of George Mackey. (Wilce, 2021)

Moreover, from time to time, new approaches to the foundations of QM in the last decades claimed to pay attention to this reaxiomatizing attempt. Just to mention a view that is quite popular in our times, the so-called relational approach to QM was originally an instance: in his first paper on the subject Rovelli claims.

\[10\text{ The reference is to two highly influential works of George W. Mackey, a 1957 paper and a 1963 book (Mackey 1957, 1963).}\]
that quantum mechanics will cease to look puzzling only when we will be able to *derive* the formalism of the theory from a set of simple physical assertions («postulates», «principles») about the world. Therefore, we should not try to append a reasonable interpretation to the quantum mechanics formalism, but rather to derive the formalism from a set of experimentally motivated postulates.” (Rovelli, 1996, p. 1639).\(^{11}\)

Whatever the historical antecedents of its ‘reaxiomatising’ attitude, the IT-view ascribes to a structure like $L(H)$ a highly relevant role, since it is a non-Boolean structure “in which there are built-in, structural probabilistic constraints on correlations between events” (Bub & Pitowsky, 2010, p. 439). This is exactly the way in which the IT-view of QM justifies a peculiar use of the kinematics/dynamics distinction:

The structure of Hilbert space imposes kinematic (i.e., pre-dynamic) objective probabilistic constraints on events to which a quantum dynamics of matter and fields is required to conform, through its symmetries, just as the structure of Minkowski space-time imposes kinematic constraints on events to which a relativistic dynamics is required to conform. In this sense Hilbert space provides the kinematic framework for the physics of an indeterministic universe, just as Minkowski space-time provides the kinematic framework for the physics of a non-Newtonian, relativistic universe. There is no deeper explanation for the quantum phenomena of interference and entanglement than that provided by the structure of Hilbert space, just as there is no deeper explanation for the relativistic phenomena of Lorentz contraction and time dilation than that provided by the structure of Minkowski space-time (Bub & Pitowsky, 2010, p. 439).

But what are the presuppositions of the use of such distinction in the quantum realm?

The kinematics/dynamics distinction, to the extent to which it is supposed to play a conceptual role in the Bub-Pitowsky approach, is explicitly inspired to a distinction firstly proposed by Einstein in a short but influential text, written in 1919 for *The Times*, the distinction between *constructive* theories and *principle* theories:

We can distinguish various kinds of theories in physics. Most of them are constructive. They attempt to build up a picture of the more complex phenomena out of the materials of a relatively simple formal scheme from which they start out. Thus the kinetic theory of gases seeks to reduce mechanical, thermal, and diffusional processes to movements of molecules—i.e., to build them up out of the hypothesis of molecular motion. When we say that we have succeeded in understanding a group of natural processes, we invariably mean that a constructive theory has been found which covers the processes in question. Along with this most important class of theories there exists a second, which I will

\(^{11}\) A more recent and developed instance is Höhn & Wever 2017.
call “principle-theories.” These employ the analytic, not the synthetic, method. The elements which form their basis and starting-point are not hypothetically constructed but empirically discovered ones, general characteristics of natural processes, principles that give rise to mathematically formulated criteria which the separate processes or the theoretical representations of them have to satisfy. Thus the science of thermodynamics seeks by analytical means to deduce necessary conditions, which separate events have to satisfy, from the universally experienced fact that perpetual motion is impossible. The advantages of the constructive theory are completeness, adaptability, and clearness, those of the principle theory are logical perfection and security of the foundations. The theory of relativity belongs to the latter. (Einstein [1919] 1954, p. 228)

Constructive theories provide a model of phenomena that is supposed to account for such phenomena in terms of a structure which is more simple and fundamental: the canonical example is the kinetic theory, able to explain thermodynamic phenomena in terms of the particles’ motion. Principle theories, on the other hand, are developed by the formulation of (empirically well-founded) generalizations, such that the theories express formal conditions that the phenomena under scrutiny are held to satisfy: according to Einstein, relativistic theories belong to this second class.12

The kinematical/dynamical distinction, in the interpretation that the Bub and Pitowsky take to be relevant in the quantum realm, is actually a variant (and not simply an extension) of the Einstein constructive/principle distinction, developed in three steps. First, Bub and Pitowsky map the kinematical/dynamical distinction onto the Einsteinian principle/constructive distinction, by pairing kinematical with principle and dynamical with constructive. Second, the <kinematical-principle> / <dynamical-constructive> [KP/DC, from now on] distinction is then articulated by applying it entirely within special relativity – whereas the original Einsteinian distinction assigns special relativity theory as a whole to the class of principle theories. In this articulation, Bub and Pitowsky adopt the context and the language of a more recent debate on the role of the KP/DC distinction in special relativity, a debate in which two diverging views confront each other on the meaning of the physical explanation of phenomena provided by special relativity. According to one view, special relativity does its job in prescribing the structure of space–time via the specification of essentially kinematic constraints, that phenomena displacing in spacetime are held to satisfy:

That a free particle moves in a straight line is kinematical in this reckoning since such trajectories are the geodesics associated with the flat affine structure of Minkowski space–time. […] I would say that Minkowski space–time

12 According to a standard reading of the Einstein presentation of the distinction, only constructive theories are really explanatory, a reading that Einstein himself appears to suggest and endorse. Since relativistic theories are included in the class of principle theories, the above reading raises the puzzling question of what might Einstein really mean when he suggests that relativistic theories do not explain the phenomena they cover: an possible way out is proposed by Lange 2014. By a historical point of view, the background of the Einstein distinction is interestingly explored in Giovanelli 2020.
encodes the default spatio-temporal behavior of all physical systems in a world in accordance with the laws of special relativity. Special relativity is completely agnostic about what inhabits or—to phrase it more awkwardly but in a way that may be more congenial to a relationist—carries Minkowski space–time. All the theory has to say about systems inhabiting/carrying Minkowski space–time is that their spatio-temporal behavior must be in accordance with the rules it encodes. Special relativity thus imposes the kinematical constraint that all dynamical laws must be Lorentz invariant.” (Janssen, 2009, pp. 27-8).

According to an alternative view, spacetime theories must receive a dynamical-constructive understanding in that their geometrical properties and structure must be shown to depend on the details of a (quantum) theory of matter:

Relativistic phenomena like length contraction and time dilation are in the last analysis the result of structural properties of the quantum theory of matter. […] one is committed to the idea that Lorentz contraction is the result of a structural property of the forces responsible for the microstructure of matter. (Brown, 2005, pp. vii-viii, 132).
In our view, the appropriate structure is Minkowski geometry precisely because the laws of physics, including those to be appealed to in the dynamical explanation of length contraction, are Lorentz covariant […] From our perspective […] it is the Lorentz covariance of the laws that underwrites the fact that the geometry of space-time is Minkowskian” (Brown & Pooley, 2006, pp. 10, 14).

On the foundations of special relativity Bub and Pitowsky side with the kinematical reading, a reading that in their view justifies the establishment of an order of relevance between kinematics and dynamics: the latter, in fact, comes first but in a sort of provisional status until a kinematical account is provided; the main historical piece of evidence that is ordinarily offered for the endorsement of this order is taken to be exactly the transition from the Lorentzian approach – a constructive one – to the Einsteinian approach – whose kinematical character is almost universally considered to be the reason why physics has glorified Einstein rather than Lorentz.13
This position leads naturally to the third and final step: namely an application of the KP/DC distinction to QM in which – so it is claimed – QM should be interpreted primarily in a principle-kinematical perspective, in which the theoretical constraints concern information transfer. This information-theoretic view would justify in turn the direct adoption of a Hilbert space structure as a fact, with the consequence that non-distributivity of this structure turns out to be a ‘kinematical’ constraint in itself:

The information-theoretic interpretation is the proposal to take Hilbert space as the kinematic framework for the physics of an indeterministic universe, just

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13 The claim that the Lorentz-to-Einstein transition should be considered an evidence for the constructive-to-kinematical transition is disputed in Acuna 2014.
as Minkowski space provides the kinematic framework for the physics of a non-Newtonian, relativistic universe.” (Bub, 2020, p. 200).

According to their supporters, this option would carry with itself for free the possibility to account for two peculiar and puzzling aspects of QM: its apparent non-locality and its apparent irreducibly probabilistic status. In fact, the non-Booleanity of the Hilbert space as-a-kinematical-framework would account both for the existence of entangled states – those states that generate peculiar ‘non-local’ correlations – and for the intrinsically non-classical probabilistic structure – since a probability theory defined over a non-Boolean structure cannot be classical:

In special relativity, the geometry of Minkowski space imposes spatiotemporal constraints on events to which the relativistic dynamics is required to conform. In quantum mechanics, the non-Boolean projective geometry of Hilbert space imposes objective kinematic (i.e., pre-dynamic) probabilistic constraints on correlations between events to which a quantum dynamics of matter and fields is required to conform. In this non-Boolean theory, new sorts of non-local probabilistic correlations are possible for ‘entangled’ quantum states of separated systems, where the correlated events are intrinsically random, not merely apparently randomlike coin tosses.” (Bub, 2020, p. 201).

As a final result, therefore, in the Bub-Pitowsky information-theoretic view it is the very non-Booleanity of the Hilbert space as-a-kinematical-framework that explains “how measurements can have definite outcomes”, (dis)solving thereby what they call the big measurement problem: if we take into account the Hilbert space structure within the IT-view of QM, since what matters in this view are only the events-as-measurement-outcomes, it follows that the structure of events-as-measurement-outcomes is best hosted by the non-commutative algebraic structure of projectors, without any need to tell a ‘dynamical’ or ‘mechanical’ (as Bub sometimes calls it) story about how such events come about.

4 The Bub-Pitowsky (dis)solution does not work

This proposal is problematic, though. As we just have seen, it is a certain application of the <kinematical-principle> / <dynamical-constructive> distinction to the quantum realm in the Bub-Pitowsky approach that motivates the claim according to which the very adoption of a mathematical structure like a Hilbert space explains away the (big) measurement problem. This application, however, rests on a specific view of what the KP/DC distinction does in special relativity: it looks safe to say that, according to this view, it is the Minkowski geometry qua mathematical structure that does the job. It is a very strong assumption, indeed, that cannot go easily undisputed. First of all, it is far from clear how is this job supposed to be exactly possible or, in slightly more precise terms, how is the mathematical structure supposed to provide constraints for the physical content of the theory. As Brown aptly pointed out,
as matter of logic alone, if one postulates space-time structure as a self-standing, autonomous element in one’s theory, it need have no constraining role on the form of the laws governing the rest of content of the theory’s models. So how is its influence on these laws supposed to work?” (Brown, 2005, p. 143).14

Moreover, in what is now the ordinary presentation of the kinematic vs. constructive debate concerning special relativity (the presentation we referred to above when introducing the Bub-Pitowsky proposal), a radical alternative is usually presented in terms of explanatory priority. Either Lorentz invariance as a physical principle explains the role of the Minkowski geometry, or the other way round. As a consequence, in each of the two camps one competitor becomes the explanans and the other becomes the explanandum: tertium non datur. Does the alternative really need to be so radical? Is it really well-posed in these aut-aut terms? In fact, in relatively recent times, plausible motivations to doubt it have been put forward. According to Pablo Acuna, for instance,

a more nuanced, adequate and fruitful construal of the explanatory foundations of special relativity is that Lorentz invariance and Minkowski structure do not constitute two features of the theory such that one has to be explained by the other. Rather, they can be understood as two sides of a single coin, so there is no need to demand for an arrow of explanation connecting them. (Acuna, 2016, p. 8, my emphasis)

According to the point of view advocated by Acuna, the “demand for an arrow of explanation” is far from mandatory: taking this demand to be mandatory is what turns both the kinematical (Janssen) and the dynamical (Brown) views into “over-interpretations of the connection between Lorentz invariance and Minkowski spatiotemporal structure” (Acuna, 2016, p. 9). In place of a strict choice between the two views, Acuna proposes to read the Minkowski spatiotemporal structure as the ‘conceptual unfolding’ of the physical content expressed by the constraint coded into Lorentz invariance. There appears to be in fact a sort of mutual implication between the geometrical and the physical dimension of special relativity as a whole theory, a mutual implication that accords well also the historical process in the development of special relativity:

what Minkowski did was not to provide for the physical grounds of the results of Einstein’s (1905) paper. Minkowski’s contribution is a conceptual display of Einstein’s work, in the sense of an overt description of the spatiotemporal structure underlying the theory—a structure that Einstein had already glimpsed (Acuna, 2016, p. 9, emphasis in the original).

This new reading of the kinematic/dynamic interplay in the foundations of special relativity is illustrated with the aid of a figure in which the role of the above mentioned mutual implication between Lorentz invariance and Minkowski spatiotemporal structure is suitably emphasized:

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14 See on this point also the discussion in Acuna 2014, Sect. 3.
I argue that this highly plausible picture casts serious doubts on the alleged (dis)solution of the measurement problem by the IT-view of QM. In this view, as we have seen, it is the non-Booleanity of the Hilbert space as-a-kinematical-framework that is supposed to explain away the (big) measurement problem, with the alleged additional advantage of happily justifying non-classical probability theory and correlations. The possibility for this Hilbert space structure to achieve all these results qua kinematical framework, however, depends exactly on reading the confrontation between Lorentz invariance on one side and Minkowski geometry on the other in terms of an ‘explanatory victory’ of the latter over the former. But if the need for this explanatory asymmetry is questioned, by taking seriously that the relation between Lorentz invariance and Minkowski geometry can be read in terms of mutual implication rather than in terms of explanatory priority of one over the other, we may also argue that the non-Booleanity of the Hilbert space as-a-kinematical-framework cannot be such a cheap (dis)solution of the measurement problem.

No matter what its merits or limitations can be, the 2010 Bub-Pitowsky formulation appears to be an especially crisp and well-characterized stance on the measurement problem within an information-theoretic framework for QM, and its use of the Einsteinian distinction < principle/constructive > is far from being a drawback in itself, since the explanatory background of the original distinction is made to operate in a different context in a stimulating way. Although this formulation is still regarded as an influential view, still debated in different directions,15 however, it is not the last word of Bub’s information-theoretic approach to QM on the measurement problem, that has undergone an evolution.16 The first step has been realized in the 2016 book Bananaworld: Quantum Mechanics for Primates, where Bub apparently does not adopt anymore the dynamics-constructive/kinematics-principle distinction, as a tool.

15 See, for instance, Allori 2020 and Ben-Menahem 2020. Moreover, one of the leading experts in the information-theoretic approach to the foundations of QM – Lucas Dunlap – on his personal website makes the case that the Bub-Pitowsky 2010 work has attracted most of the analysis on “whether Quantum Information Theory offers the interpretational resources to effectively deal with the important questions—the measurement problem in particular—raised by the proponents of the psi-ontic interpretations of quantum theory.” (www. http://lucasdunlap.com/research/, accessed 9/1/2023). Timpson 2010 and Dunlap 2015 had already provided critical evaluations of the 2010 Bub-Pitowsky approach, albeit with arguments that are significantly different from the argument developed in the present paper.

16 I am grateful to a reviewer for pushing me to emphasize this evolution.
that may serve for properly locating such problem(s) within the foundations of quantum mechanics itself. He still assumes, however, that the non-commutative Hilbert space structure – taken as the natural ‘realization’ for the structure of information – plays a ‘kinematic’ role in the very same spirit of the dynamics/kinematics distinction. In fact, this assumption continues to be entirely motivated by the analogy with special relativity, and the role of the non-commutative Hilbert space structure is completely analogous to the role that (Bub argues) is played by the Minkowski geometry in dictating the structure of possible events in space–time.

On the information-theoretic interpretation, the relation between quantum mechanics and the structure of information is analogous to the relation between special relativity and the structure of space-time. Heisenberg’s “re-interpretation” of classical quantities as noncommutative, or more specifically the intertwining of commuting and noncommuting observables, imposes objective pre-dynamic probabilistic constraints on correlations between events in a similar sense to how Minkowski space-time imposes kinematic constraints on events. […] On the information-theoretic interpretation, the “big” measurement problem is a pseudo-problem. […] The random selection of a definite outcome is a feature of the nonclassical structure of information in a quantum world, just as Lorentz contraction is a feature of the Minkowski structure of space-time in a relativistic world. (Bub, 2016, p. 223).

By this point of view, therefore, the critical argument formulated in the present paper still applies to the 2016 version of the Bub information-theoretic (dis)‘solution’ of the measurement problem in quantum mechanics: even the 2016 version inherits the petitio principii according to which positing a certain abstract mathematical structure (Minkowski geometry in special relativity, Hilbert space structure in quantum mechanics) has an explanatory value in itself, no matter what the interplay of such structure with their intended phenomena is supposed to be. A still different twist to Bub’s information-theoretic view of quantum mechanics can be found in the 2018 paperback edition of his 2016 book. In order to see what this twist amounts to, we can start from a hint included again in the 2010 Bub-Pitowsky paper. At the page 451 we read:

Applying quantum mechanics kinematically, say in assigning probabilities to the possible outcomes of a measurement of some observable of a microsystem, we consider the Hilbert space of the relevant degrees of freedom of the microsystem and treat the measuring instrument as simply selecting a Boolean subalgebra in the non-Boolean event space of the microsystem to which the Born probabilities apply.” (emphasis added).

In a way, the Bub, 2018 takes seriously this hint: the move is then to shift the focus from the underlying non-Boolean event space, whose ontological status is
left unexplained (or simply taken to need no further explanation), to a reading of the measurement process in essentially Bohrian terms, in which the choice of a (quasi-transcendental) measurement context and the need for an effective linguistic communication of the measurement outcomes have a formal counterpart in the ‘selection’ of a Boolean subalgebra, out of the overall non-Boolean event structure. Bub explicitly concedes that his current information-theoretic interpretation “has more in common with what I think Bohr and some other of the early proponents of the Copenhagen interpretation had in mind” (Bub, 2018, p. 227) and, in this direction, aptly quotes a 1939 essay in which Bohr claims that.

the only significant point is that in each case some ultimate measuring instruments [...] must always be described entirely on classical lines, and consequently be kept outside the system subject to quantum mechanical treatment” (Bohr, 1998, p. 104, my emphasis).

It is then entirely consistent on Bub’s part to emphasize a crucial feature of this revised view of the measurement within his most recent formulation, namely its perspectival character: “On the information-theoretic interpretation of Hilbert space, just one observer perspective is legitimate in the application of quantum mechanics: the perspective of the observer for whom an actual measurement occurs at the macrolevel.” (Bub, 2018, p. 227). As a consequence, Bub may well claim that

the information-theoretic interpretation rejects […] the assumption that quantum mechanics is universal – not by restricting the universality of the dynamics or any part of quantum mechanics, but by placing a constraint on how the theory is applied. (Bub, 2018, p. 228)

Now, the extent to which this ‘return to Bohr’ is actually an advancement over the Bub-Pitowsky 2010 formulation is debatable. On the one hand, the overall information-theoretic framework is ‘lighter’, since any explicit or implicit suggestion that this framework describes information as a sort of basic stuff of the universe is definitively dropped (Dunlap, 2022). But, on the other hand, in its latest version it appears to inherit entirely the problematic status of a Bohrian interpretation of the measurement process: in this interpretation, quantum mechanics simply has no clearly defined resources to fix what John S. Bell called “the shifty split” between the quantum world (in Bub’s terms, the world of the non-Boolean structure of events) and the classical world (in Bub’s terms, the Boolean subalgebra selected by the choice of a measurement context). In the apt formulation of David Mermin:

Bell also deplored a “shifty split” that haunts quantum mechanics. The shiftiness applies both to the nature of the split and to where it resides. The split can be between the quantum and the classical, the microscopic and the macroscopic, the reversible and the irreversible, the unspeakable (which requires the quantum formalism for its expression) and the speakable (which can be said in ordinary language). In all cases the boundary is moveable in either direction, up to an ill-defined point. Regardless of what is split from what, all versions of the shifty split are vague and ambiguous. (Mermin, 2012, p. 8).
Moreover, along with this return to Bohr, the assumption of the kinematic role of the non-Boolean structure as a fact is retained also in Bub, 2018, as the following passage clearly shows:

I use the term ‘information-theoretic’ to emphasize that the quantum revolution is about new sorts of probabilistic correlations in nature, analogous to the sense in which the relativistic revolution is about new sorts of spatio-temporal relations. A theory of information is fundamentally a theory of probabilistic correlations, and that’s what Hilbert space gives you. Probabilities and probabilistic correlations arise as a feature of the non-Boolean structure. They are “uniquely given from the start”, to quote von Neumann, not measures over states as they are in a classical or Boolean theory. The interwinement of commuting and noncommuting observables in Hilbert space imposes objective pre-dynamic probabilistic constraints on correlations between events, analogous to the way in which Minkowski space-time imposes kinematic constraints on events. (Bub, 2018, pp. 224-225).18

In closing this section, let me return to the mutual implication view. If we take it to be a plausible view for special relativity, we may wonder whether this view can carried over the quantum domain, by asking whether a kind of mutual implication between physical principles on one side and the Hilbert space structure on the other can be imagined along the lines of the pair Lorentz invariance/Minkowski geometry: in this case, the Hilbert space structure (with its formal constraints, such as the non-distributivity of the structure of projection operators) might be operating as the ‘unfolding’ of physical principles that do tell a ‘mechanical’ story – in Bub’s terms – about how results come about in the measurement process. The plausibility in principle of a mutual-implication account for QM raises at least two points. The first is that, unlike the case with the IT-view in which we remain essentially within the boundaries of standard Copenhagenish QM, we have to select an interpretation that is able in principle to tell such kind of story (that is, either Bohmian mechanics or GRW). This is totally resonant with the fact that this species of interpretations take seriously the existence of a genuine measurement problem in QM, but still this move forces us to confront with the details of these interpretive frameworks. The second is that, in the case of special relativity, the physical side – Lorentz invariance – is represented by a sort of meta-nomological statement, and there seems to be nothing in the quantum realm that plays a similar role at the level of the foundations of the theory.

No matter how serious this last difference should be considered, let us take into account a physical principle (referred to as ‘Doctrine Q’19) that might play in the quantum world a role that is somewhat similar to the role that Lorentz invariance

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18 In this sense, I agree with the Dunlap evaluation: “while the claim that [the information-theoretic view of quantum mechanics] is a realist interpretation of quantum mechanics is deemphasized in Bananaworld, almost all of the structure of the 2010 view remains, including the dissolution of the measurement problem.” (Dunlap 2022, p. 44).

19 This is the terminology used by Lazarovici, Oldofredi, Esfeld 2018 to denote this principle, to which they refer in the form given by Mermin 1993, p. 803: “It is a fundamental quantum doctrine that a measurement does not, in general, reveal a preexisting value of the measured property.”
plays for the kinematical/dynamical distinction regarding special relativity. A possible formulation of this doctrine goes as follows (Goldstein et al., 2011):

**Doctrine Q** – *It is a general principle of quantum theory that measurements of physical quantities do not simply reveal pre-existing or pre-determined values, the way they do in classical theories. Instead, the particular outcome of the measurement somehow "emerges" from the dynamical interaction of the system being measured with the measuring device.*

In support of the claim that the mutual-implication account might be significant for my argument against the IT-view treatment of the measurement problem, let us consider then the ‘Doctrine Q’ in a specific interpretational setting, Bohmian mechanics, in order to see how in this case the Hilbert space mathematical machinery turns out to be a suitable theoretical environment for the ‘unfolding’ of the content of such a physical principle.

Standard Bohmian mechanics (BM from now on) is an observer-free formulation of non-relativistic QM that does not endorse the completeness axiom of standard QM, according to which the wave function encodes the maximal amount of information that is possible to extract concerning the state of the physical quantum systems. This stance is implemented by adding to the wave function the information concerning the system *position*: as a consequence, the wave function – in addition to satisfying the Schrödinger equation – determines the particles’ motion via the especially Bohmian addition to the ordinary structure of quantum mechanics, namely the so-called guiding equation. Therefore, standard BM describes quantum particles and their trajectories in physical space and time: in doing this, BM is said to provide a space–time ontology of non-relativistic quantum mechanics, namely a class of well-specified kind of objects and properties displayed in space–time that quantum mechanics is supposed to be about (being the primary target of the theory, this space–time ontology is called *primitive ontology*). In slightly more precise terms, the main assumptions of BM are (Lazarovici et al., 2018):

*Particle configuration:* There always is a configuration of $N$ permanent point particles in the universe, characterized only by their positions $X_1, \ldots, X_N$ in three-dimensional, physical space at any time $t$.

*Guiding equation:* A wave function $\Psi$ is ascribed to the particle configuration that, at the fundamental level, is the universal wave function attributed to all the particles in the universe together.

*Schrödinger equation:* The wave function always evolves according to the Schrödinger equation.

*Typicality measure:* On the basis of the universal wave function $\Psi$, a unique ‘typicality measure’ can be defined in terms of the $|\Psi|^2$–density.\(^{20}\) Given that typicality measure, it can then be shown that for nearly all initial conditions, the distribution of particle configurations in an ensemble of sub-systems of the universe that admit of a wave function $\psi$ of their own (known as *effective wave function*) is

\(^{20}\) For a proof of the uniqueness result, see Goldstein & Struyve 2007.
a $|\psi|^2$–distribution. A universe in which this distribution of the particles in sub-configurations obtains is considered to be in quantum equilibrium.

Two consequences of this theoretical framework are especially important for our purposes. First, under the assumption of quantum equilibrium for the Bohmian universe, the Born rule for the calculation of measurement outcome statistics on sub-systems of the Bohmian universe can be derived (where the ordinary $\psi$ for particular sub-systems within the universe is their effective wave function). Second, BM does not introduce intrinsic properties for the (sub)systems it is about except position, since the effective wave function that describes is provided at any time $t$ by a pair $(X_t, \psi_t)$, where $X_t$ describes the actual spatial configuration of the system. These two consequences jointly illustrate a possible sense in which the Hilbert space structure might operate as a suitable abstract environment for the ‘unfolding’ of such a physical principle, and this is due to the especially relevant, and twofold, role that the wave function plays in that structure. By the physical point of view, the wave function grounds the computation of probability for measurement outcomes and prevents the attribution of pre-existing properties to systems subject to measurement; in turn, both these aspects are somehow ‘mirrored’ in the mathematical status of the wave function itself within the Hilbert space structure (sanctioned in particular by the Gleason theorem): it is in this sense that a sort of mutual implication between the physical content of Doctrine Q and the mathematical language of quantum theory might be seen at work. The circumstance that, within Bohmian mechanics, this hypothetical mutual implication is preserved – since the above two-fold role and status of the wave function can be derived from the main assumption of the interpretation – together with the simplicity and non-ambiguity of the overall Bohmian ontological picture of the quantum world, might provide support for conceiving the latter as a suitable framework for a consistent reading of the Doctrine Q-Hilbert space relationship.

5 The In principle underdetermination claim

In the above pages, a major role has been played by the confrontation between the ‘kinematical’ virtues of the Minkowskian Einstein theory and the ‘dynamical’ drawbacks of the Lorentz theory, where the contrived, conspiratorial and ad hoc character of the latter was especially relevant to convince the majority about its unplausibility. In his 2005, Bub attacks Bohmian mechanics by arguing that it plays with respect to the IT-view of QM the role that the Lorentz theory played with respect to the Einstein one: this attack is based on what Henderson called a form of in principle underdetermination claim (Henderson, 2020). In the IT-view of QM, the assumption of the relevant information-theoretic constraints is sufficient – according to the CBH theorem – to single out a Hilbert space-based theoretical structure for quantum phenomena: since Bohmian mechanics is predictively equivalent to Hilbert space QM by construction, Bub argues that
a constructive theory like Bohm’s theory can have no excess empirical content over a quantum theory. Just as in the case of Lorentz’s theory, Bohm’s theory will have to posit contingent assumptions to hide the additional mechanical structures (the hidden variables will have to remain hidden), so that in principle, as a matter of physical law, there could not be any evidence favouring the theory over quantum theory. (Bub, 2005, pp. 555)

The argument, relying on the impossibility of empirically discriminating between standard QM and Bohmian mechanics (this is what in principle undetermination is supposed to mean), clearly suggests that the motivations underlying the adoption of Bohmian mechanics, as an alternative interpretation of standard QM, should be considered as suspicious as the motivations underlying the adoption of the Lorentz theory with respect to special-relativistic phenomena. In other words, Bohmian mechanics vis-a-vis standard QM would as contrived as Lorentz theory vis-a-vis Einstein theory only because the former ‘can have no excess empirical content’ over the latter. Moreover, in Bub’s view, the same fate occurs to the GRW model: although the latter may differ in principle from standard QM in terms of empirical predictions, both Bohmian mechanics and the GRW model are taken to needlessly ‘add structure’ to standard QM and therefore are subject to the same argument:

On the information-theoretic interpretation, no assumption is made about the fundamental ‘stuff’ of the universe. So, one might ask, what do tigers supervene on? In the case of Bohm’s theory or the GRW theory, the answer is relatively straightforward: tigers supervene on particle configurations in the case of Bohm’s theory, and on mass density or ‘flashes’ in the case of the GRW theory, depending on whether one adopts the GRWm version or the GRWf version. […] The solutions to the ‘big’ measurement problem provided by Bohm’s theory and the GRW theory are dynamical and involve adding structure to quantum mechanics. There is a sense in which adding structure to the theory to solve the measurement problem dynamically—insofar as the problem arises from a failure to recognize the significance of Hilbert space as the kinematic framework for the physics of an indeterministic universe—is like Lorentz’s attempt to explain relativistic length contraction dynamically, taking the Newtonian spacetime structure as the underlying kinematics and invoking the ether as an additional structure for the propagation of electromagnetic effects. In this sense, Bohm’s theory and the GRW theory are ‘Lorentzian’ interpretations of quantum mechanics. (Bub & Pitowsky, 2010, pp. 452, 454)

The argument of Bub and Pitowsky is far from convincing. As far as Bohmian mechanics is concerned, the focus on the absence of excess empirical content as the exclusive criterion for comparing the pairs Lorentz/Einstein and Bohmian mechanics/QM fails to distinguish between the particular ways in which the Lorentz theory on one side and Bohmian mechanics on the other fare concerning the issue of the empirical indistinguishability from the respective rival theory. Whereas in the Lorentz theory a form of in-principle undetectability of physical effects is
introduced in a totally unconventional and instrumental way, focused as it is to
the very preservation of ether as the privileged frame, the overall aim of Bohmian
mechanics is to describe quantum phenomena as much as possible in line with the
long and honored tradition of physics in which the theory is supposed to be about
‘matter in motion’. Moreover, a major factor in the interpretation is represented
by a scientifically respectable account of inaccessibility of the particles’ position:
although the position of every quantum system is definite at all times, we are de
facto unable to control each such position, an uncontrollability which should look
far from surprising in a quantum world anyway, and which after all resembles other
forms of inaccessibility to which we are used also in a pre-quantum world, such as
in statistical classical mechanics. In this respect, the absence of ‘excess empirical
content’ might be taken to be a price that we are willing to pay if we can have in
return a view of the microphysical world as a more familiar world of particles in
motion. There seems to be no apriori argument why this sort of explanatory virtues
should be considered less valuable than the alleged explanatory virtue of informa-
tion-theoretic constraints, a virtue that manifests itself in the one and only effect of
being the basis for deriving a Hilbert space structure for quantum phenomena.

Moreover, there is a further point. If the measurement problem is “the problem
of explaining how individual measurement outcomes come about dynamically”, a
solution at hand is already available on the market, since the account of the meas-
urement problem (ordinarily understood) provided by the dynamical reduction
approach of GRW does exactly that. In principle, the GRW approach proposes a
dynamical account of the measurement interaction in terms of a rigorous recipe
on how – in the event of a measurement of a physical quantity on a system S by an
apparatus A – the (non-linearly modified) dynamics of the joint system S+A physi-
cally produces a definite result out of the initial, entangled state of S+A:

This approach consists in accepting that the dynamical equation of the stand-
ard theory should be modified by the addition of stochastic and nonlinear
terms. The nice fact is that the resulting theory is capable, on the basis of a
single dynamics which is assumed to govern all natural processes, to account
at the same time for all well-established facts about microscopic systems as
described by the standard theory, as well as for the so-called postulate of wave
packet reduction (WPR), which accompanies the interaction of a microscopic
system with a measuring device.” (Bassi & Ghirardi, 2020).

Due to the non-linear modification of the ordinary quantum dynamics, the GRW
model does have in principle an excess empirical content, an evidence that might in
principle favour the GRW model over standard QM. Since it is possible to articu-
late an ontological reading of the GRW model, it appears highly controversial to
claim that “no mechanical theory of quantum phenomena that includes an account
of measurement interaction can be acceptable.” (Bub, 2004, p. 241), although spec-
ifying the sense in which the GRW model can tell a ‘mechanical story’, accord-
ing to the different, possible underlying ontologies, is a non-trivial matter (Bassi &
Ghirardi, 2020).
6 Conclusions

At least since the times of the Wittgensteinian stance displayed in the *Tractatus*, an attitude toward open foundational issues proved especially tempting: that of dissolving rather than solving the problem at stake, by showing that under certain assumptions there is simply nothing to solve, namely that it is the very status of problem that need not hold for the claim under scrutiny. This is the option that Jeffrey Bub and Itamar Pitowsky developed about the infamous measurement problem in non-relativistic QM, in the framework of their information-theoretic view of quantum theory. In the present paper I have tried to show that their ‘deconstructing’ strategy is far from convincing. Not because of the central role that the notion of information plays for the foundations of the theory – this is a totally legitimate view, among the many present in the wide market of the interpretations of quantum theory – but rather because the strategy rests entirely on the claim that an exclusively ‘kinematical’ reading of special relativity constitutes a good account of explanation for the phenomena covered by this theory: so good that it can be extended to QM and exploited, in order to show that the (exclusively ‘kinematical’ reading of the) Hilbert space structure in itself is analogously a good explanation for the quantum facts, so that we do not need any detailed, ‘dynamical’ or ‘mechanical’ account of the measurement process from which the quantum facts themselves emerge. I emphasized that if there can be grounds to question the very necessity for the explanatory priority of a kinematical account over a dynamical one (or vice versa), then there can be grounds correspondingly for rejecting what I take to be a cheap (dis)solution of the measurement problem in QM.

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Declarations

Ethical approval Not Applicable.

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Conflict of interest The author declares that he has no conflict of interest.

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