Electrochemical Fabrication of rGO-embedded Ag-TiO₂ Nanoring/Nanotube Arrays for Plasmonic Solar Water Splitting

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HIGHLIGHTS

• Reduced graphene oxide (rGO) in electrode can weaken the light scattering of plasmonic Ag nanoparticles and promote the hot electrons transfer from Ag nanoparticles to Ti substrate.

• A route synergizing rGO with plasmonic Ag on TiO₂ for plasmonic solar water splitting was provided.

ABSTRACT  Effective utilization of hot electrons generated from the decay of surface plasmon resonance in metal nanoparticles is conductive to improve solar water splitting efficiency. Herein, Ag nanoparticles and reduced graphene oxide (rGO) co-decorated hierarchical TiO₂ nanoring/nanotube arrays (TiO₂ R/T) were facilely fabricated by using two-step electrochemical anodization, electrodeposition, and photoreduction methods. Comparative studies were conducted to elucidate the effects of rGO and Ag on the morphology, photoresponse, charge transfer, and photoelectric properties of TiO₂. Firstly, scanning electron microscope images confirm that the Ag nanoparticles adhered on TiO₂ R/T and TiO₂ R/T-rGO have similar diameter of 20 nm except for TiO₂ R-rGO/T. Then, the UV–Vis DRS and scatter spectra reveal that the optical property of the Ag-TiO₂ R/T-rGO ternary composite is enhanced, ascribing to the visible light absorption of plasmonic Ag nanoparticles and the weakening effect of rGO on light scattering. Meanwhile, intensity-modulated photocurrent spectroscopy and photoluminescence spectra demonstrate that rGO can promote the hot electrons transfer from Ag nanoparticles to Ti substrate, reducing the photogenerated electron–hole recombination. Finally, Ag-TiO₂ R/T-rGO photoanode exhibits high photocurrent density (0.98 mA cm⁻²) and photovoltage (0.90 V), and the stable H₂ evolution rate of 413 μL h⁻¹ cm⁻² within 1.5 h under AM 1.5 which exceeds by 1.30 times than that of pristine TiO₂ R/T. In line with the above results, this work provides a reliable route synergizing rGO with plasmonic metal nanoparticles for photocatalysis, in which, rGO presents a broad absorption spectrum and effective photogenerated electrons transfer.

KEYWORDS  TiO₂ nanoring/nanotube hierarchical structure; Reduced graphene oxide; Spectral responses; Plasmonic Ag nanoparticles; Water splitting
1 Introduction

Nowadays, more and more researchers have paid attention to the preparation and application of nanomaterials. Strong photosresponse in the visible light region allows plasmonic metal nanoparticles to serve as tunable light antennas of wide-bandgap semiconductors, such as TiO$_2$, for overcoming the limitations of rapid recombination of photogenerated electron–hole pairs and inability to absorb visible light [1]. Among the plasmonic metals, Ag nanocrystals have attracted considerable attention for their remarkable plasmon resonance enhancement. Lian et al. [2] designed a plasmonic Ag/TiO$_2$ photocatalytic composite by selecting Ag quantum dots to act as the photosensitizer for driving the visible light-driven photoelectrocatalytic hydrogen evolution. Chen et al. [3] deposited Ag nanoparticles on the flame reduced TiO$_2$, which exhibits significant application prospects for the enhanced solar conversion efficiency. Mohammadi et al. loaded Ag nanoparticles on TiO$_2$ nanotubes by sequential chemical bath deposition, and the photoelectrochemical (PEC) activity is considerably higher than the bare TiO$_2$ (about 3 times), owing to the light absorption improvement of plasmonic effects in addition to better separation and transport of electron–hole pairs [4]. In addition, our research group has also probed plasmonic Ag nanoparticles on TiO$_2$ nanotube arrays electrode for efficient solar water splitting [5]. Once Ag nanoparticles (Ag NPs) get in touch with TiO$_2$, a Schottky barrier is formed and the height plays an important role in electron transfer [6]. Under the ultraviolet irradiation, the electrons belonging to the valence band of TiO$_2$ migrate to its conduction band and then are transferred to Ag NPs with the help of the Schottky barrier. Under the visible light irradiation, Ag NPs exhibit surface plasmon resonance (SPR) effect, and a number of hot electrons are generated via Landau damping [7]. The high-energy hot electrons in Ag NPs can be transferred across the Schottky barrier to TiO$_2$, followed by photocatalytic reaction [8, 9]. However, the short lifetime of hot electrons (<160 fs) has been the main limiting factor of its effective utilization [10]. To make full use of plasmon-induced hot electrons, it is necessary to optimize the structure and composition of Ag/TiO$_2$ [11].

Graphene, an extensively used two-dimensional nonmetallic conductive material, is considered as an ideal electron transfer medium in photocatalysis for its various advantages, such as good interfacial contact with adsorbents, excellent mobility of charge carriers, and large surface area [12]. As a derivative of graphene, reduced graphene oxide (rGO) has the similar structure and properties [13, 14]. Lang et al. [15] made rGO nanosheets as conductive “bridge” to help transferring the hot electrons from Ag of smaller work function to TiO$_2$ of larger work function. Gao et al. [16] fabricated the composite photocatalyst with Ag and TiO$_2$ nanoparticles wrapped by rGO sheets, in which, rGO was used as an electron acceptor. In most designs, both Ag and TiO$_2$ are co-deposited on the surface of rGO nanosheets to form Ag/rGO/TiO$_2$ ternary catalysts. The two-dimensional rGO nanosheets act as a good electron transporting bridge owing to its high electron mobility and π-electron conjugation [17, 18]. Besides, rGO can broaden the absorption spectrum of TiO$_2$ [19]. For example, a well-organized rGO and Ag wrapped TiO$_2$ nanohybrid was successfully achieved by Leong and his colleagues, and the wrapped rGO nanosheets promoted visible light shift toward red spectrum [20].

The Ag/rGO/TiO$_2$ ternary nanocomposite is usually used in heterogeneous systems. For instance, Ong et al. developed hybrid organic PVDF-inorganic Ag-rGO-TiO$_2$ nanocomposites for multifunctional volatile organic compound sensing and photocatalytic reaction [21]. Nasrollahzadeh prepared Ag/rGO/TiO$_2$ nanocomposite for the reduction of 4-nitrophenol, Congo red and methylene blue in aqueous media [22]. Pant et al. synthesized a multifunctional Ag-TiO$_2$/rGO powder for wastewater treatment [23]. Recently, Tian et al. [24] synthesized mesoporous TiO$_2$/rGO/Ag by an electrostatic self-assembly approach and a photo-assisted reduction process for the degradation of methylene blue. As much as we know, there is little research on utilizing the Ag/rGO/TiO$_2$ ternary catalysts in PEC system. In addition, there are so many methods to fabricate various conformations of Ag/rGO/TiO$_2$ ternary composite, such as hydrothermal synthesis for 3D urchin-like Ag/TiO$_2$/rGO composites [25], electrodeposited and chemical reduction methods for Ag/rGO co-decorated TiO$_2$ nanotube arrays [26, 27], electrospinning technique for TiO$_2$ nanofibers loaded on rGO/Ag platform [28], three-step method for TiO$_2$-Ag-rGO vertical heterostructure [29], and so on, but the effect of rGO on the other two composites has been rarely discussed in these researches.

In the present work, TiO$_2$ nanoring/nanotube arrays (TiO$_2$ R/T) via two-step anodization were chosen as substrate electrode, which shows unique oscillating absorption in the visible region and excellent photoelectric properties in our
previous work [30]. Herein, we focus on fabricating the TiO$_2$
R/T electrodes with Ag NPs and rGO fragments to further
improve the H$_2$ production rate in PEC system. Different
from other studies, rGO was added into the electrolyte of the
first or second electrochemical anodization. Thus, the effects
of rGO on morphology, photoresponse, charge transfer, and
photoelectric properties of TiO$_2$ and Ag in the nanocom-
posites were systematically investigated. Accordingly, the
possible charge transfer routes and light absorption of as-
prepared samples were discussed.

2 Methods

2.1 Materials

Ti foils (0.25 mm thick, 99.5% purity) were supplied by Alfa-
Aesar Company. Aluminum oxide (Al$_2$O$_3$) was purchased
from Gaona Powder Company (Shanghai, China). Graphite
oxide powders were purchased from Xianfeng nanomaterials
Company (Nanjing, China). Silver nitrate (AgNO$_3$), sodium
dodecyl sulfate (SDS), polyvinylpyrrolidone (PVP), and
sodium borohydride (NaBH$_4$) were purchased from Fucheng
Chemical Reagent Company (Tianjin, China). Ethylene
glycol (EG), sodium nitrate (NaNO$_3$), ammonium fluoride
(NH$_4$F), sodium sulfate (Na$_2$SO$_4$), triethanolamine (TEOA),
ethanol (C$_2$H$_6$O) and acetone (CH$_3$COCH$_3$), isopropyl alco-
hol (C$_3$H$_8$O) were purchased from Beijing Chemical Works
(Beijing, China). All reagents were of analytical grade and
used without further purification.

2.2 Preparation of TiO$_2$ R/T Modified with rGO
in Nanoring or Nanotube

TiO$_2$ R/T was prepared by using two-step anodization
method in ethylene glycol solution with NH$_4$F (0.25 wt%)
and water (2 vol%) [30]. The first and second anodic oxi-
dation steps correspond to the formation of nanorings and
nanotubes, respectively. The graphene oxide (GO) solution
of 0.75 g L$^{-1}$ was prepared by adding given amount of GO
powder to deionized water and sonicated for 5 h. A certain
amount of SDS was added to enhance exfoliation and sep-
erate graphite oxide sheets from each other. Subsequently,
0.6970 g NH$_4$F and 5 mL above GO-dispersed solution were
dissolved in 250 mL ethylene glycol. The mixed solution
was stirred uniformly. In the first anodization, Ti foil was
anodized at 60 V for 60 min using ethylene glycol solution
with NH$_4$F and H$_2$O as electrolyte, and the as-grown nano-
tube layer was ultrasonically removed in deionized water.
After that, the pretreated Ti foil was used as anode again
for the second anodization in the ethylene glycol solution
with NH$_4$F and GO-dispersed solution. Finally, the sample
was reduced in 1 mol L$^{-1}$ NaBH$_4$ aqueous solution for 1 h
with 300 W Xe-lamp irradiation and subsequently dried at
130°C for 2 h in air. In the progress, the oxygen-containing
functional groups on the surface of GO were removed and
the $sp^2$ hybrid orbits of carbon were re-created [31, 32]. The
prepared sample is named as TiO$_2$ R/T-rGO. The other sam-
ple, prepared in NH$_4$F with rGO for the first anodization
and NH$_4$F with H$_2$O for the second anodization, is named
as TiO$_2$ R-rGO/T.

2.3 Deposition of Ag Nanoparticles

A two-electrode setup was used for pulse electrodепosition
of Ag NPs by using TiO$_2$ R/T, TiO$_2$ R/T-rGO, and TiO$_2$
R-rGO/T samples as the working electrodes, and Pt sheet
as the counter electrode. The current pulsing approach was
utilized via choosing a cathodic pulse (− 25 mA, 0.1 s) and
a relation time (0.3 s) in a mixed solution containing 1 mM
AgNO$_3$, 10 mM NaNO$_3$, 0.2 g L$^{-1}$ PVP, and 0.8 vol% iso-
propyl alcohol with magnetic string at 200 rpm. The as-obtained
products are labeled as Ag-TiO$_2$ R/T, Ag-TiO$_2$ R/T-rGO,
and Ag-TiO$_2$ R-rGO/T. The preparation conditions and cor-
responding samples are summarized and listed in Table 1.

2.4 Characterizations

The morphologies of TiO$_2$ R/T hierarchical structures
were characterized with a field-emission scanning electron
microscope (FESEM; Hitachi, Japan, S4800). Transmis-
sion electron microscopies (TEM) were obtained with a
JEM-2100 transmission electron microscope (Japan Elec-
tron Optics Laboratory Co., Ltd., JEOL) at an accelerat-
ing voltage of 200 kV. The Raman spectra were acquired
using a confocal micro-Raman system (BRUKER, RFS100)
with an excitation wavelength of 532 nm. The atomic force
microscopic (AFM) measurement was performed with a
custom-designed scanning probe microscope (NTegra
Spectra, NTMDT), using a Co/Cr-coated silicon cantilever.
The crystal phases of the samples were determined using an X-ray diffractometer (XRD; D8 Advance Bruker/AXS) with Co (30 mA) radiation. The diffraction data were recorded for 2θ between 5° and 80°. X-Ray photoelectron spectroscopy (XPS) (Escalab 250Xi, Thermo Fisher Co., USA) measurements were processed using an Al-Kα monochromatic X-ray source (1486.6 eV). All binding energies were referenced to the C 1 s peak at 284.8 eV of surface adventitious carbon. The optical absorption spectra were obtained by employing a Fiber Spectral Instrument (AVANTES, Netherlands) and compressed BaSO4 powder as reflectance standard. Photoluminescence (PL) studies were done using a Gilden Photonics photoluminescence spectrophotometer under 260 nm excitation. Scattering spectra were obtained under 410–590 nm excitation.

2.5 Photoelectrochemical Measurements

The PEC performances were measured in a three-electrode system with TiO2-based composites, Ag/AgCl electrode, and Pt mesh as the working, reference, and counter electrodes, respectively. The supporting electrolyte was carried out in aqueous solution containing 15 vol% TEOA and 0.5 M Na2SO4. The samples were illuminated by a solar simulator (91160, Newport) which equipped with an AM 1.5 filter. The intensity was calibrated to 100 mW cm−2. The visible light activity was evaluated under chopped light irradiation with light wavelength ≥ 400 nm by using a filter (FSQ-GG400, Newport). The linear sweep voltammetry (LSV), photocurrent density–time (I-t), open-circuit photovoltage (U-t), and intensity-modulated photocurrent spectroscopy (IMPS) results were collected by an electrochemical workstation (AUTOLAB PGSTAT302 N). The wavelength of the irradiation light for IMPS measurements was set as 410 nm. Detection of hydrogen was performed online by using an Agilent 7890B gas-chromatograph (5Å molecular sieve, N2 carrier gas).

2.6 Computational Model and Methods

The general gradient approximate (GGA) with the PBE was adopted to estimate the optical properties of graphene using CASTEP code [33]. An energy cutoff of 750 eV has been used for expanding the Kohn–Sham wave functions. The SCF tolerance was set as 2.0 × 10−6 eV/atom. The k-points grid sampling of Monkhorst–Pack scheme was set as 1 × 1 × 1 in the irreducible Brillouin zone. Moreover, the fast Fourier transform grid was set as 48 × 48 × 48. The dipole corrections were utilized for the model, which are essential in eliminating the nonphysical electrostatic interaction between periodic images.

3 Results and Discussion

3.1 The Effect of rGO on the Compositions and Morphologies of Ag/TiO2

The compositions of bare and modified TiO2 R/T were obtained by XRD patterns, as shown in Fig. 1a. For all samples, the diffraction peaks could be indexed to the tetragonal TiO2 anatase phase (JCPDS No. 21-1272) and the hexagonal Ti metal phase (JCPDS No. 44-1294). rGO and Ag have little effect on the crystal phase of TiO2 R/T. However, no rGO diffraction peaks could be identified due to its high dispersity and low loading, which may be below the detection limit of XRD [27]. The diffraction peaks ascribed to Ag NPs were also not clearly distinguished because the peak corresponding to Ag at 44.6° could be covered up by the peak of Ti at 44.9° [34]. Therefore, the XPS technique was employed to

| Table 1 Preparation conditions of the TiO2-based samples |
|---------------------------------------------------------|
| Samples | The electrolyte using in the first anodization | The electrolyte using in the second anodization | Deposition of Ag NPs |
| TiO2 R/T | NH4F + H2O + (CH2OH)2 | NH4F + H2O + (CH2OH)2 | No |
| Ag-TiO2 R/T | NH4F + H2O + (CH2OH)2 | NH4F + H2O + (CH2OH)2 | Yes |
| TiO2 R/T-rGO | NH4F + H2O + (CH2OH)2 | NH4F + GO solution + (CH2OH)2 | No |
| Ag-TiO2 R/T-rGO | NH4F + H2O + (CH2OH)2 | NH4F + GO solution + (CH2OH)2 | Yes |
| TiO2 R-rGO/T | NH4F + GO solution + (CH2OH)2 | NH4F + H2O + (CH2OH)2 | No |
| Ag-TiO2 R-rGO/T | NH4F + GO solution + (CH2OH)2 | NH4F + H2O + (CH2OH)2 | Yes |
analyze the specific surface composition and elemental binding energy, as displayed in Fig. 1b. It is suggested that the composites contain C, Ag, Ti, and O, and the chemical binding energies located at 284.8, 368.7, 459.5, and 529.6 eV are assigned to the characteristic peaks of C 1s, Ag 3d, Ti 2p, and O 1s, respectively. Moreover, Fig. 1c exhibits the high-resolution spectrum of Ag 3d from the Ag-modified samples. The Ag 3d_{5/2} and Ag 3d_{3/2} core levels of Ag-TiO_2 R/T, Ag-TiO_2 R/T-rGO, and Ag-TiO_2 R-rGO/T were fitted with two peaks at the binding energy of about 367 and 373 eV, respectively, corresponding to the Ag^{+} ions produced by the electron transfer between metallic Ag and TiO_2/rGO [16,29]. It should be noted that the characteristic peaks of the Ag 3d from Ag-TiO_2 R/T-rGO and Ag-TiO_2 R-rGO/T have positively shifted by 0.4 and 0.3 eV relative to Ag-TiO_2 R/T, respectively, confirming that the electrons indeed transfer from Ag to rGO. Coincidentally, Shi et al. reported that the two peaks of Ag nanotriangle both had a positive shift of 1.1 and 1.0 eV when mixed with graphene. They explained it as the electron transfer from Ag to graphene [12]. Wu et al. [35] also deduced the Ag 3d peaks displaying nearly 0.6 eV shift to higher binding energy for Ag deposited on graphene sheet is the result of electron transfer from metallic Ag to graphene sheet. With 3 nm etching, a significant positive shift of the binding energy for Ag 3d_{5/2} relative to about 368.4 eV of the bulk Ag^{+} was identified. The above results confirm that the existing state of Ag NPs is related to TiO_2 and rGO.

TEM and AFM techniques were used to test the structure and size of GO. The TEM image of GO is shown in Fig. 2a, and the ultrathin wrinkles confirm the existence of GO sheets. Through analyzing the GO flakes with AFM (Fig. 2b), the slice sizes were measured in the range of 0–200 nm and the thickness of the thin layer was about 0.9 nm. The reduction of GO in the composites was verified by Raman spectra, as displayed in Fig. 2c, exhibiting the characteristic D and G bands of the carbon materials. The intensity ratio of D to G band ($I_D/I_G$) can be used to evaluate rGO graphitization degree because it can be increased by the chemical reduction of GO [36]. The result shows that the $I_D/I_G$ ratio increases from 0.94 of GO to 1.18 of rGO.

Fig. 1  a XRD patterns, b XPS full spectrum scanning, and c high-resolution XPS spectra of the Ag 3d peaks of TiO_2-based composites
indicating the successful reduction of GO under UV irradiation [25, 37].

The as-prepared GO dispersion was dividedly put into the electrolyte of the first and second step anodic oxidation to prepare TiO$_2$ R-rGO/T and TiO$_2$ R/T-rGO, causing a few differences in $I$-$t$ and $U$-$t$ of the preparation process. The variation of current versus time is recorded in Fig. 3a; the measured current significantly gets large with GO in electrolyte, such as the current of TiO$_2$ R-rGO/T in the first step (green line) and TiO$_2$ R/T-rGO in the second step (blue line). The abundant negatively charged functional groups of GO make it moved toward the anode under the action of external potential, facilitating the formation of the TiO$_2$ hierarchical structures. On this basis, Ag NPs were deposited on TiO$_2$ by pulsed-electrodeposition method. The pulse currents were set as $-0.25$ mA of the nucleation processes and 0 mA of the growth processes, and the corresponding $U$-$t$ variations are shown in Fig. 3b. In contrast, the potential in the whole process corresponding to Ag-TiO$_2$ R/T is the smallest ($-5.55$ V, $-1.88$ V). It can be concluded that the resistance of the TiO$_2$ R/T sample in nucleation stage is the largest and its internal electric field at growth stage is the strongest, which is conducive to grow rather than nucleate of Ag NPs. In other words, the decoration of rGO in TiO$_2$ nanorings and nanotubes has reduced the internal resistance of the composites.

SEM images can clearly reveal the morphologies of TiO$_2$ R/T, Ag-TiO$_2$ R/T, TiO$_2$ R/T-rGO, Ag-TiO$_2$ R/T-rGO, TiO$_2$ R-rGO/T, and Ag-TiO$_2$ R-rGO/T. From Fig. 4a, it can be seen that the hierarchical top-ring/bottom-tube structures were successful fabricated. The average diameter of the nanorings is about 150 nm with 35 nm thickness, and the

![Fig. 2](image-url) a TEM and b AFM images of the GO flakes with the thickness distribution shown in the inset. c Raman spectra of GO and rGO

![Fig. 3](image-url) a The $I$-$t$ curves in the two-step anodic oxidation process and b $U$-$t$ curves in pulse electrodeposition of TiO$_2$ hierarchical structure
diameter of nanotubes is about 50 nm. The concaves in the top layer are expected to work as nanomirrors for light reflection and scattering. The insert gives a cross-sectional view, showing that the nanotubes with a length of about 9 μm are perpendicular to the Ti substrate. Figure 4c, e shows the TiO₂ R/T decorated with rGO in nanotubes and nanorings, respectively. It is seen that the addition of rGO has no apparent effect on the morphology of TiO₂ R/T, but the electrostatic force between rGO and Ti⁴⁺ made the nanotubes fairly grown up in the first anodization of TiO₂ R-rGO/T, thereby taking a long time to remove the film of nanotubes by ultrasound, which is consistent with Fig. 3a [38]. Due to the slow penetration of precursor solution into the air-filled nanotubes, plasmonic Ag NPs were mainly loaded on the nanorings of TiO₂ R/T [39]. In Fig. 4b, d, the average size of Ag NPs adhered on the top edge of the nanorings is about 20 nm, possessing an optimal quality factor of SPR [40]. Figure 4f shows that the average diameter of Ag NPs deposited on TiO₂ R-rGO/T has decreased to 6 nm, due to the strong Ag-C bond between Ag and rGO existing in the surface [29]. The electrostatic force is conducive to the uniform formation of Ag NPs on the surface, i.e., promotes the nucleation process. The EDS mapping over Ag-TiO₂ R-rGO/T was further investigated and presented in Fig. 4g, indicating that all the elements (O, Ti, Ag, and C) were uniformly dispersed.

3.2 The Effect of rGO on the Optical Properties of TiO₂-Based Composites

Figure 5 shows the UV–Vis absorption spectra of the as-prepared samples. Obviously, all samples show oscillating absorption peaks in the visible region, which are derived from the interference of lights reflecting from the top nanorings and the bottom Ti substrate [30]. It is obvious that the light absorption of TiO₂ R/T, TiO₂ R/T-rGO, TiO₂ R-rGO/T in the range of 350–565 nm is weak, as shown in Fig. 5a. From Fig. 5b, the deposition of Ag has greatly improved the photoabsorption of TiO₂ in the region, in agreement with the absorption enhancement of 400–650 nm in Ref. [31], which is related to the SPR effect of spatially confined electrons in Ag NPs. It has been reported that the amplitude of oscillating peaks of pristine TiO₂ R/T increases significantly with similar frequency after annealing, attributed to the formation of defects in TiO₂ [30]. Here, the absorption peaks of rGO-decorated samples show a slight red shift and the amplitude of oscillation peaks has decreased. Thus, we can infer that the addition of rGO into TiO₂ nanorings or nanotubes has changed the spatial shapes or interactions of the molecules by introducing defects into TiO₂, influencing the transition of the valence electrons. Finally, the UV–Vis absorption spectra of TiO₂ R/T with or without rGO show different performances. In addition, compared with bare TiO₂ R/T, TiO₂ R/T-rGO has no significant change, while TiO₂ R-rGO/T has weakened absorption in 300–350 nm. The same thing happens to Ag-TiO₂ R/T, Ag-TiO₂ R/T-rGO, Ag-TiO₂ R-rGO/T. It is suggested that adding rGO into nanotubes in the second step is a great way to improve the light absorption of TiO₂ R/T. Taking all samples into account, Ag-TiO₂ R/T-rGO has the greatest UV–Vis absorption property.

For the spherical metal nanoparticles, the absorption increases with the decrease in scattering [41]. At present, the main controversy regarding the mechanism of the emission is the assignment to either radiative recombination of hot carriers (photoluminescence) or light scattering [42, 43]. Here, we measured the scattering spectra of the TiO₂-based samples. A broad excitation wavelength range (410–590 nm) was chosen to contain the extinction spectrum of Ag and rGO as far as possible [44]. The final scattering spectra were collected and illustrated in Fig. 6. With the incident wavelength increasing, the scattering intensities of all samples decreased first and then increased, and the maximum scattering intensity corresponded to 900 nm irradiation within the scope of the measurement. In comparison with TiO₂ R/T, the scattering intensity of Ag-TiO₂ R/T is enhanced, as shown in Fig. 6a, b. It is known that Ag NPs are helpful to the radiative scattering process induced by coupling to a propagating plasmon mode, so the fluorescence enhancement is observed [45]. For the utilization efficiency of plasmon-excited hot carriers, one of the major limitations is the loss in carrier energy due to carrier scattering within the plasmonic metals [46]. When adding rGO into the electrodes, the drastically weakened scattering of TiO₂ R/T-rGO and Ag-TiO₂ R/T-rGO is observed in Fig. 6c, d. On the other hand, compared with TiO₂ R/T, the scattering intensity of TiO₂ R-rGO/T corresponding to incident light λ > 530 nm is higher (Fig. 6e), but the scattering intensity of
Ag-TiO$_2$ R-rGO/T has decreased to 570 (Fig. 6f), indicated by the blue dashed line. Anyway, Ag-TiO$_2$ R/T-rGO has the weakest scattering intensity, corresponding to the great light absorption as shown in Fig. 5.

To further study the effect of rGO on the optical properties of TiO$_2$-based composites, the investigations of rGO using the plane wave ultrasoft pseudopotential method based on the density functional theory were conducted in detail. The Fermi energy is set to zero, represented by the blue dashed line. The computed band structure of rGO (Fig. 7a) matches well with the previous results in Ref. [47], indicating that the character of density of states (DOS) and optical properties are reasonable and reliable. Figure 7b shows the DOS of graphene, suggesting that the electron density in Fermi energy level is the lowest. The band structures of the majority spin and the minority spin were obtained by spin calculation (Fig. 7c). The remarkable difference is that the Fermi level intersects with the majority spin while lying in the band gap of the minority spin. Therefore, graphene has apparent spin polarization and should be regarded as “half metal.” In Fig. 7d, the absorption coefficients of the anisotropic material under polarized and unpolarized light irradiation have a little difference. The absorption range of graphene is 1.64–50.00 eV, corresponding to 756.10–24.80 nm, which is calculated by Eq. 1,

$$E = \frac{hc}{\lambda}$$

where $E$ is the energy of light (eV), $h$ is Planck constant ($4.13566743 \times 10^{-15}$ eV s), $c$ is the speed of light ($3.153 \times 10^{17}$ nm s$^{-1}$), and $\lambda$ is the wavelength of light (nm).
Accordingly, rGO can effectively enhance the light absorption of TiO$_2$.

### 3.3 Photoelectrochemical Properties and Hydrogen Production Rates

Chronoamperometric I-t curves are recorded in Fig. 8a by irradiating the electrodes with AM 1.5, which are measured to investigate the transient photocurrent response. Upon illumination, the photocurrents of all photoanodes can rapidly reach a steady state and revert to zero when light turns off. All samples exhibit fast photocurrent response and stable photocurrent density under chopped light. The maximum photocurrent density is 0.98 mA cm$^{-2}$ of Ag-TiO$_2$ R/T-rGO, which is almost 1.5 times of bare TiO$_2$ R/T (0.69 mA cm$^{-2}$). The photocurrents of TiO$_2$ R/T-rGO (0.80 mA cm$^{-2}$) and Ag-TiO$_2$ R/T (0.73 mA cm$^{-2}$) are also greater than TiO$_2$ R/T, but TiO$_2$ R-rGO/T (0.63 mA cm$^{-2}$) and Ag-TiO$_2$ R-rGO/T (0.51 mA cm$^{-2}$) have smaller photocurrent density.

To characterize the photocurrent density as a function of applied voltage, a set of LSV curves are presented in Fig. 8b, measured with the applied potential window of -1.5 to 1.0 V. The obtained photocurrent densities at 0 V of the considered six samples lie in the following order: Ag-TiO$_2$ R/T-rGO > TiO$_2$ R/T-rGO > Ag-TiO$_2$ R/T > TiO$_2$ R-rGO/T > Ag-TiO$_2$ R-rGO/T. It is consistent with the measured photocurrent densities under AM 1.5 shown in Fig. 8a. In addition, the apparent oxidation peaks of Ag-TiO$_2$ R/T, Ag-TiO$_2$ R/T-rGO, and Ag-TiO$_2$ R-rGO/T at about 0.25 V correspond to the transformation of Ag$^0$ into Ag$^+$. Figure 8c shows the open-circuit photovoltage curves, suggesting the n-type conductivity of all samples [48]. The negative increase in voltage under light irradiation suggested that the photogenerated electrons are injected from the semiconductor film into the Ti substrate [49]. The differences of the potential in light and dark are calculated as 0.68, 0.84, 0.65, 0.90, 0.65, and 0.74 V, corresponding to TiO$_2$ R/T, Ag-TiO$_2$ R/T, TiO$_2$ R/T-rGO, Ag-TiO$_2$ R/T-rGO, TiO$_2$ R-rGO/T, and Ag-TiO$_2$ R-rGO/T, respectively. The open-circuit potential has increased with the loading of Ag NPs, while a slight decrease of about 0.03 V has been observed with the addition of rGO in nanotubes and nanorings. All above results show that the strategy of designing Ag-TiO$_2$ R/T-rGO nanostructure is efficient for promoting the photoelectric properties of the PEC system. Combined with the images shown in Fig. 4d, f, the Ag NPs deposited on TiO$_2$ R/T-rGO and TiO$_2$ R-rGO/T are 20 and 6 nm, respectively. Therefore, although rGO can reduce the size of Ag NPs, the TiO$_2$ composite electrode deposited by Ag NPs with the strongest SPR effect has better photoelectric performance.

In accordance with the UV–Vis absorption spectra (Fig. 5), the enormous enhancement of photocurrent density under visible light (λ > 400 nm) illumination is observed in Fig. 8d. Obviously, Ag-loaded samples give quick response and then get down to the stable values. The interconversion of Ag/Ag$^+$ during the alternating irradiation ensures a stable visible light-induced photocurrent [50]. The maximum photocurrent density is 0.06 mA cm$^{-2}$ of Ag-TiO$_2$.
R/T-rGO, followed by 0.017 mA cm\(^{-2}\) of Ag-TiO\(_2\) R/T, 0.015 mA cm\(^{-2}\) of TiO\(_2\) R/T, and 0.005 mA cm\(^{-2}\) of TiO\(_2\) R/T-rGO. Thus, Ag-TiO\(_2\) R/T-rGO hybrid exhibits a fourfold increased current density under visible light compared with pristine TiO\(_2\) R/T.

At last, we tested the hydrogen generation performance of Ag NPs and rGO enhanced hierarchical TiO\(_2\) nanoring/nanotube arrays via choosing pristine TiO\(_2\) R/T for comparison. The electrolyte was deaerated with N\(_2\) for 3–4 h until the dissolved oxygen could be neglected. Once the tested three-electrode system exposed on simulated sunlight, the rate of H\(_2\) evolution rapidly increased during the first hours and remained nearly constant afterward, as described in Ref. \[51\]. The responding time of TiO\(_2\) R/T and Ag-TiO\(_2\) R/T-rGO is 60 and 120 min, respectively. As displayed in Fig. 9, Ag-TiO\(_2\) R/T-rGO electrode exhibits excellent photoelectrochemical water splitting activity under AM 1.5, with an enhanced average H\(_2\) evolution rate of 413 \(\mu\)L h\(^{-1}\) cm\(^{-2}\). This value exceeds by 1.30 times than that of TiO\(_2\) R/T (317 \(\mu\)L h\(^{-1}\) cm\(^{-2}\)). It remains stable during the test, partly due to the suitable conduction band edge of rGO, which is high enough to make H\(_2\) generated \[52, 53\].

3.4 The Charge Transfer and Proposed Mechanisms of TiO\(_2\)-Based Electrodes with Decoration of Ag and rGO

Intensity-modulated photocurrent spectroscopy (IMPS) is usually used to examine the recombination of conduction band electrons at surface states and charge transfer to solution. Here, we used IMPS to clarify the relationship of the charge transfer with SPR effect of Ag as well as rGO. From the IMPS complex plane plots, the electron-transport time \((\tau_d)\) can be calculated with the frequency \((f_{\text{min}})\) at the zenith of the semicircle through Eq. 2,

\[
\tau_d = \frac{1}{2\pi f_{\text{min}}}
\]

As shown in Fig. 10a, there is one distinct semicircle for TiO\(_2\) R/T and TiO\(_2\) R/T-rGO, and their electron-transport time is calculated as 6.34 and 4.00 ms, respectively. Ag-TiO\(_2\) R/T and Ag-TiO\(_2\) R/T-rGO exhibit two semicircles, implying two different electron-transport modes \[54\]. As for Ag-TiO\(_2\) R/T, a fraction of the electrons were collected with a short transport time (0.20 ms), while the remaining electrons were collected with a long time (20 ms). The electron-transport time of Ag-TiO\(_2\) R/T-rGO also contains
two parts, and the total time equals to 16.42 ms, which is shorter than that of Ag-TiO$_2$ R/T. Thus, rGO can accelerate the charge transfer over TiO$_2$ nanotube arrays. It is helpful for the migration of plasmon-induced hot electrons from Ag to TiO$_2$.

By analyzing the photoluminescence (PL) spectra of the TiO$_2$-based composites, the lifetime of the photogenerated electron–hole pairs in semiconductors under irradiation can be obtained. The greater the photoluminescence emission, the more recombination of electron–hole pairs. Thus, the Ag-TiO$_2$ R/T-rGO sample exhibits lower electron–hole recombination rate, as revealed by its weakened PL peak (Fig. 10b). At the same time, Ag-TiO$_2$ R/T and TiO$_2$ R/T-rGO also show a lower recombination rate than TiO$_2$ R/T, owing to the trapping of the photo-excited electrons of Ag NPs and the high charge carrier mobility of rGO, respectively [17].

On the basis of the discussion above, the photoresponse and charge transfer play an important role in PEC hydrogen production via water splitting. Therefore, the proposed mechanisms of Ag-TiO$_2$ R/T-rGO are schematically illustrated in Fig. 11. In this picture, the gray sphere and hexagon represent Ag NPs and rGO flakes, respectively. On the one hand, UV light (purple arrows), which can be absorbed by TiO$_2$, and visible light (red arrows) interfere with each other in the nanotubes producing oscillation peaks, as shown in Fig. 5. Part of the photogenerated electrons is transferred to Ag NPs to reduce the loss of holes in recombination process. Meanwhile, the light irradiated on rGO is absorbed (turn yellow). On the other hand, the light
in the range of 350–565 nm (green arrows) is absorbed by Ag NPs (turn yellow) through the SPR excitation, and then, the hot electrons are injected into the conduction band of TiO$_2$ ($E_{cd} = -4.2$ eV) over the Schottky barrier [5, 9]. The wall of TiO$_2$ nanotubes possesses excellent ability to transfer electrons, allowing more electrons to be trapped by H$^+$ with the formation of H$_2$ [36]. Moreover, the surface of the Ag NPs has inevitably oxidized to Ag$_2$O ($E_{cd} = -4.47$ eV, $E_{vd} = -5.98$ eV), according to Fig. 1c. Because the energy of the electrons in the conduction band is insufficient to reduce a proton, Ag$_2$O serves as an electrocatalyst for the oxidation reaction of photoelectrical water splitting [55, 56]. Eventually, both of the photogenerated electrons and hot electrons are transported through Ti substrate to Pt photocathode to perform the hydrogen evolution reaction, and yet holes are enriched synchronously on Ag$_2$O to oxidize water into O$_2$.

Fig. 8  a Photocurrent responses in the light on–off process under illumination of AM 1.5, b Linear sweep voltammetry, c open-circuit photovoltage curves, and d photocurrent responses in the light on–off process under illumination of visible light ($\lambda > 400$ nm) of TiO$_2$-based nanomaterials

Fig. 9 Photocatalytic H$_2$ production. The hydrogen evolution rate was measured by using the TiO$_2$ R/T and Ag-TiO$_2$ R/T-rGO samples as photoanodes and Pt foil as cathode in a PEC cell containing a 0.5 M Na$_2$SO$_4$ solution under AM 1.5
Fig. 10  a IMPS plots and b PL spectra with the excited wavelength at 260 nm of TiO$_2$ R/T, Ag-TiO$_2$ R/T, TiO$_2$ R/T-rGO, and Ag-TiO$_2$ R/T-rGO

Fig. 11  Schematic illustration of the photoresponse and photogenerated charge transfer of Ag-TiO$_2$ R/T-rGO under UV–Vis light irradiation
4 Conclusions

In summary, we designed the hybrid Ag-TiO$_2$-rGO catalyst structure to realize enhanced wide spectrum photocatalytic hydrogen evolution and mainly researched the great influence of rGO on Ag and TiO$_2$. For pristine TiO$_2$ R/T and TiO$_2$ R/T with rGO in nanotubes, 20 nm Ag NPs were uniformly deposited on the surface. When adding rGO into TiO$_2$ nanorings, the size of the deposited Ag NPs has reduced to 6 nm due to the Ag-C bond between Ag and rGO. With the ternary composites formed, the electron transfer from Ag to rGO was identified by the positively shifted of the Ag 3d peaks in XPS spectra. The deposition of Ag NPs on TiO$_2$ R/T can enhance the absorption in 350–565 nm region, but the light scattering of Ag-decorated samples was also enhanced attributed to the propagating plasmon mode. Fortunately, the introduction of rGO can greatly weaken the scattering due to its broad absorption in UV–Vis light region, verified by the density functional theory calculations. On the other hand, IMPS illustrates that rGO can promote the transportation of hot electrons to Ti substrate, resulting in the reduction of photogenerated electron–hole recombination, as shown in PL plots. By comparison, Ag-TiO$_2$ R/T-rGO photoelectrode shows the largest photocurrent density and open-circuit potential under AM 1.5 irradiation. Meanwhile, the remarkable photocurrent under visible light irradiation highlights its excellent performance in visible region. Finally, Ag-TiO$_2$ R/T-rGO photoelectrode resulted in a remarkable boost in the H$_2$ evolution rate (413 μL h$^{-1}$ cm$^{-2}$) compared to pristine TiO$_2$ R/T photoelectrode (317 μL h$^{-1}$ cm$^{-2}$). Collectively, our work opens a new window for the novel design and synthesis of rGO-contained plasmonic photoelectrode, which is conducive to give full play to the excellent performance of rGO.

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