Hierarchical Star Formation across the ring galaxy NGC 6503

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ABSTRACT

We present a detailed clustering analysis of the young stellar population across the star-forming ring galaxy NGC 6503, based on the deep HST photometry obtained with the Legacy ExtraGalactic UV Survey (LEGUS). We apply a contour-based map analysis technique and identify in the stellar surface density map 244 distinct star-forming structures at various levels of significance. These stellar complexes are found to be organized in a hierarchical fashion with 95% being members of three dominant super-structures located along the star-forming ring. The size distribution of the identified structures and the correlation between their radii and numbers of stellar members show power-law behaviors, as expected from scale-free processes. The self-similar distribution of young stars is further quantified from their autocorrelation function, with a fractal dimension of ~1.7 for length-scales between ~20 pc and 2.5 kpc. The young stellar radial distribution sets the extent of the star-forming ring at radial distances between 1 and 2.5 kpc. About 60% of the young stars belong to the detected stellar structures, while the remaining stars are distributed among the complexes, still inside the ring of the galaxy. The analysis of the time-dependent clustering of young populations shows a significant change from a more clustered to a more distributed behavior in a time-scale of ~60 Myr. The observed hierarchy in stellar clustering is consistent with star formation being regulated by turbulence across the ring. The rotational velocity difference between the edges of the ring suggests shear as the driving mechanism for this process. Our findings reveal the interesting case of an inner ring forming stars in a hierarchical fashion.

Key words: galaxies: spiral – stars: formation – galaxies: stellar content – galaxies: individual (NGC 6503) – galaxies: structure – methods: statistical
1 INTRODUCTION

Star formation, the conversion of gas into stars, is a key process in shaping the structure, morphology and evolution of galaxies. Young star clusters, OB associations and large complexes of young stars are the signposts of the recent star formation across a galaxy. These various young stellar concentrations, covering a large dynamic range in physical length-scales, do not appear to be independent from each other but rather related to each other in a hierarchical fashion. Large loose structures of stars host smaller and more compact star-forming systems, which themselves are substructured (e.g., Elmegreen et al. 2000). The characterization of this clustering behavior and its origins are key issues in understanding how star formation progresses in space and time across galactic scales. Samples of resolved young stellar populations across a whole galaxy allow the detailed investigation of their formation in structures, improving our understanding from integrated stellar light.

There are only a few investigations on the clustering of resolved young stars across galactic scales. They are focused on the Magellanic Clouds (Maragoudaki et al. 1998; Gieles et al. 2008; Bastian et al. 2009), NGC 6822 (Gouliermis et al. 2010), and M 33 (Bastian et al. 2007), up to sub-Mpc distances, where individual stars could be resolved from the ground. These studies give evidence of the self-similar scaling relations in the parameters of the identified stellar structures, demonstrating the hierarchical nature of stellar structural morphology on galactic scales. With the present study we extend these investigations to larger distances in the Local Volume, based on the exquisite resolving ability of the Hubble Space Telescope (HST) with its cameras Advanced Camera for Surveys (ACS) and Wide-Field Camera 3 (WFC3). We also add to the sample of investigated galaxies the interesting case of the ring (possibly barred) spiral galaxy NGC 6503 (Fig. 1).

NGC 6503 is classified as of morphological type SAB(s)bc:1, i.e., a pure s-shaped spiral galaxy, with possibly well-developed spiral arms, showing a trace of a bar (Buta et al. 2015). The galaxy has also a patchy circumnuclear appearance in Hα, interpreted by Knapen et al. (2006) as a nuclear star-forming ring. The same authors postulate that while this galaxy is possibly one of the rare cases of ring spirals classed as unbarred2, it is most likely to actually be barred. Indeed, both the H I velocity field (Bottema & Gerritsen 1997) and near-IR imaging (Freeland et al. 2010) support the presence of a strong end-on bar. Moreover, Freeland et al. (2010) argue that the previously identified nuclear ring is instead an inner ring spanned in diameter by the bar (Fig. 1). According to theory, rings form by gas accumulation through the action of gravitational torques from the bar pattern (e.g., Simkin et al. 1980; Schwarz 1984). They are thus sites of active star formation in the galaxy (Buta & Combes 1996).

Being a ring galaxy observed by HST, NGC 6503 is a unique case for investigating in detail star formation across an inner ring. Specific questions that can be probed are: (a) What is the length-

1 This type is according to the classification by Buta et al. (2015) from the Spitzer Survey of Stellar Structure in Galaxies (S4G). The galaxy was previously classified according to the Third Reference Catalog of Bright Galaxies (de Vaucouleurs et al. 1991) as SA(s)cd, i.e., a non-barred spiral with loosely wound spiral arms, and with no apparent ring.

2 Rings are most probably resonance phenomena, caused by a rotating bar or other non-axisymmetric disturbance in the disk. The current evidence supports the idea that “rings are a natural consequence of barred galaxy dynamics” (Buta & Combes 1996).
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2 OBSERVATIONS AND PHOTOMETRY

LEGUS is a HST panchromatic stellar survey of 50 nearby star-forming dwarf and spiral galaxies with an emphasis on UV-enabled science applications. Images in a wide waveband coverage from the near-UV to the I-band are being collected with WFC3 and ACS in parallel, and combined with archival optical ACS data. The survey, its scientific objectives and the data reduction are described in Calzetti et al. (2015). Stellar photometry will be described in detail in Sabbi et al. (in prep). The images of NGC 6503 to be used in our analysis were obtained in the filters F275W, F336W, F438W, F555W and F814W (equivalent to NUV, U, B, V, and I respectively) in August 2013. We applied the pixel-based correction for charge-transfer efficiency (CTE) degradation using tools provided by STScI, before processing them with ASTRODRIZLE and prior to their photometry.

Photometry was performed with the package DOLPHOT (e.g., Dolphin 2000). This package performs point-spread function (PSF) fitting using PSFs especially tailored to HST cameras. Before performing the photometry, we first prepared the images using the DOLPHOT packages ACSMASK and SPLITGROUPS. Respectively, these two packages apply the image defect mask and then split the multi-image STScI FITS files into a single FITS file per chip. We then used the main DOLPHOT routine to make photometric measurements in each filter independently on the pre-processed images, relative to the coordinate system of the drizzled reference image. The output photometry from DOLPHOT is on the calibrated VEGA-MAG scale based on the zeropoints provided on the WFC3 page.

Indicative Hess diagrams of the complete stellar samples retrieved with this process are shown in Fig. 2. From these diagrams it is seen that different wavelengths cover different stellar types, with young stars in the blue filters and old stars in the red filters.
It should be noted that our stellar photometry may include unresolved stellar clusters or multiple systems. However, the criteria applied in the following section for the selection of the best photometric measurements eliminate this contamination to the minimum. In any case, such systems, as well as unresolved binary systems, which should be still included in the best photometric sample, have a negligible effect to our analysis, because the measured blue light of every source is determined by that of the dominant bright star in the system.

2.1 Selection of the Stellar Sample

We separate the stellar sources with the most reliable photometry into four catalogs, each including stars found in two adjacent filters throughout the complete wavelength coverage from F275W (UV) to F814W (I). The waveband coverage of the filter-pairs overlap, so that, e.g., the bluer pair (F275W, F336W) shares filter F336W with the next redder pair (F336W, F438W), which shares filter F438W with the next and so forth. Since bluer filters cover the younger populations, while the redder ones identify mostly the more evolved stars, our selection distinguishes stars at roughly different evolutionary stages. These best photometrically defined stellar samples are determined in terms of quality parameters returned by our photometry. We apply the selection of this sample for stellar sources identified successfully in at least two adjacent filters. We compile the final photometric stellar catalogs in each filter-pair by applying the following quality criteria:

- **DOLPHOT type of the source**, \( \text{TYPE} = 1 \)
- **Crowding of the source in each of the filters**, \( \text{CROWD} < 2 \)
- **Sharpness of the source squared in each filter**, \( \text{SHARP}^2 < 0.3 \)
- **Signal-to-noise ratio in each filter**, \( \text{SNR} > 5 \)

In **DOLPHOT**, the object type parameter has a value of 1 for the best stars in the photometry. Stars too faint for PSF determination and non-stellar sources have \( \text{TYPE} > 1 \). The crowding parameter is a measure of how much brighter the star would have been measured had nearby stars not been fit simultaneously. For an isolated star it has the value of zero. The sharpness is zero for a perfectly-fit star,
We compared these distributions against the light distribution from Spitzer images from Dale et al. (2009) in 4.5 μm, tracing old stars, and 8 μm, an indicator based on dust emission of where the young stars are, in order to check whether the observed stellar distributions could be affected by dust attenuation. This comparison confirmed that the blue stars are more clustered in a ring-like structure, traced by the 8 μm image, while the 4.5 μm image shows a smooth distribution for the old populations. Nevertheless, the distributions shown in Fig. 6 (top panel) suffer from projection effects that limit our analysis on the global stellar clustering through the observed stellar distributions. Thus, we correct the observed stellar positions for the known inclination of the galaxy.

### 2.2 Correction for projection of the galactic disk

In order to assess the spatial distribution of stars, free from the effect of projection of the galactic disk, we apply a deprojection of the stellar positions from the plane of the sky to the plane of the galaxy, under the simple assumption of an axisymmetric flat rotating disk for NGC 6503. This assumption is based on the remarkably regular gas kinematics in NGC 6503 that are well described by rotation only (Kuzio de Naray et al. 2012).

For a star with position (x, y) in the plane of the sky (where the center of the galaxy is at the origin), a simple rotation by the position angle, φ, yields

$$\begin{align*}
x' &= x \sin(\phi) - y \cos(\phi) \\
y' &= x \cos(\phi) + y \sin(\phi).
\end{align*}$$

(1)

The coordinates of the star in the mid-plane of the galaxy are then x’ and y’/cos(i), i being the inclination angle. We rotated our stellar catalogs for a position angle of φ = 135°, determined by the orientation of the observed field-of-view. We then corrected them for projection for an inclination angle of i = 75.1°, over its kinematic center (J2000) 17°49′26.30″, 70°08′40.7″ (both determined by Greisen et al. 2009). The spatial distributions of the blue and red stellar samples corrected for projection are shown in Fig. 6 (bottom panel). It is interesting to note that the deprojected distribution of the blue young stellar population (Fig. 6 – bottom-left panel) is in remarkable agreement of the deprojected GALEX NUV data, presented by Freeland et al. (2010). Both distributions highlight the star-forming ring of the galaxy and its nucleus LINER (low-ionization nuclear emission-line regions; Lira et al. 2007), shown as a UV-brigh clump at the center of the ring (Freeland et al. 2010, their Fig. 5). In the following section we further compare the distribution of our blue stellar sample with that of the red sample, through their stellar surface density maps.

### 2.3 Stellar Surface Density Maps

We construct stellar surface density maps for the blue and red samples with the application of the Kernel Density Estimation (KDE), by convolving the stellar catalogs with a Gaussian kernel. The FWHM of this kernel specifies the “resolution” at which stellar structures will be revealed. These density maps can be treated as significance maps, corresponding to two dimensional probability functions for the positions of the stars in each sample. For their construction with this method, the most important (and debatable) parameter is the size of the KDE kernel (the convolving

| Sample | Filter pair | Total number of sources | Number of sources not in bluer filters | Percentage over whole sample |
|--------|-------------|-------------------------|----------------------------------------|-----------------------------|
| Blue   | F275, F336  | 12834                   |                                        |                             |
| Green  | F336, F438  | 13499                   | 1570                                   | 12%                         |
| Yellow | F438, F555  | 21438                   | 8484                                   | 40%                         |
| Red    | F555, F814  | 42511                   | 27525                                  | 65%                         |

More details about the fit-quality parameters are given in DOLPHOT documentation, available at http://americano.dolphinsim.com/dolphot/

It is interesting to note that the 4.5 μm image shows only a faint signature of a bar.
FWHM), i.e., the resolution of the constructed maps. At the distance of NGC 6503 (∼ 5.3 Mpc; Karachentsev et al. 2003), 1 second of arc corresponds to ∼ 25.7 parsec, and therefore every WFC3 UVIS pixel (∼ 0.04") corresponds to a physical scale of ∼ 1 pc. Consequently a KDE map constructed from our data with a resolution of ∼ 10 pc requires a kernel with FWHM of ∼ 10 pixels (∼ 0.4 arcsec). However, such a map would be extremely noisy, revealing small compact stellar over-densities rather than coherent physically related stellar concentrations. The KDE kernel size to be applied depends on the science to be achieved from the KDE maps, and it is best decided based on experimentation. Our tests on various kernel sizes showed that a FWHM of ∼ 80 pc, comparable to the scale of typical OB associations (Gouliermis 2011, and references therein) and giant molecular clouds (GMCs; Bolatto et al. 2008, and references therein), is the most appropriate for the detection of large star-forming structures.

The stellar surface density maps of NGC 6503, constructed from our photometry with a kernel of ∼ 80 pc (80 UVIS pixels), for both the blue and red stellar samples are shown in Fig. 7. From the KDE map of the blue sample (Fig. 7 – left panel) it is seen that the young blue stellar population depicts the inner star-forming ring of NGC 6503 (e.g., Mazzuca et al. 2008; Freeland et al. 2010), where individual large structures can be identified. On the other hand, the red stellar sample found in filters F555W and F814W is much more spread out, populating the whole extent of the disk of the galaxy. In the following sections we further explore the clustering behavior of young stars in NGC 6503. Specifically, we (a) conduct a census of young stellar structures, and determine their demographics, across the extent of the observed field-of-view, and (b) investigate
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Figure 7. Surface stellar density maps constructed with the Kernel Density Estimation method with a kernel of FWHM ~ 80 pc, for stars identified with the best photometry in the filter pairs F275W, F336W (left panel), and F555W, F814W (right panel). The color bar at the top corresponds to stellar density significance levels in values of $\sigma$, where $\sigma$ is the density standard deviation in the whole of each sample. The ‘blue’ stellar population surface density tracks very well the star-forming ring of NGC 6503 (e.g., Mazzuca et al. 2008; Freeland et al. 2010). LEGUS photometry reveals individual star-forming structures across the ring in order to access the large-scale progression of star formation in space and time. On the other hand, the ‘red’ evolved stellar population is far more extensively distributed than the blue, populating the whole disk of the galaxy. This difference is consistent with the more prolonged dynamical mixing of more evolved populations, and their dissolution in the field of NGC 6503.

the global distribution of the young populations and assess its hierarchy in order to understand the global star formation topology in NGC 6503.

3 YOUNG STELLAR STRUCTURES IN NGC 6503

The blue stellar population selected in the previous sections for the study of clustered star formation comprises 12 834 stars. Distinct concentrations of these young stars can be identified as stellar over-densities in the KDE stellar surface density maps. These over-densities correspond to individual star-forming structures, such as stellar associations, stellar aggregates and stellar complexes, the large-scale centres of star formation in a galaxy. The identification and characterization of these systems is thus important to our understanding of how galaxies construct their stellar components and what is their future evolution. In this section we compile a detailed catalog of all young stellar concentrations that can be revealed from our observations of NGC 6503. Our method is applied by first constructing the stellar surface density KDE map of the blue stars. With 80 pc resolution we are able to identify loose stellar concentrations that correspond to the stellar associations and stellar complexes of the galaxy, while avoiding a high noise level produced by random small-scale density fluctuations. An iso-density contour plot of this surface density map is shown in Fig. 8 (left panel). In this plot, contour lines correspond to different density levels, starting from that corresponding to the average density ($0\sigma$) up to the highest level of $16\sigma$. They are drawn in steps of $1\sigma$ with different colors accordingly. These isopleths determine the borders of various stellar structures identified across NGC 6503 at various significance levels. The physical dimensions of each detected structure are defined by these borders. A parameter considered in our detection is the minimum number of stars included in an over-density in order to be classified as a structure, which is $N_{\text{min}} = 5$, in line with other identification techniques (e.g., Bastian et al. 2007). This criterion eliminates the detection of random stellar congregations, so-called asterisms. Additionally, we consider as real those structures which appear in at least two

3.1 Detected Young Stellar Structures

Young stellar structures are identified on the KDE stellar surface density map of the blue stars. With 80 pc resolution we are able to identify loose stellar concentrations that correspond to the stellar associations and stellar complexes of the galaxy, while avoiding a high noise level produced by random small-scale density fluctuations. An iso-density contour plot of this surface density map is shown in Fig. 8 (left panel). In this plot, contour lines correspond to different density levels, starting from that corresponding to the average density ($0\sigma$) up to the highest level of $16\sigma$. They are drawn in steps of $1\sigma$ with different colors accordingly. These isopleths determine the borders of various stellar structures identified across NGC 6503 at various significance levels. The physical dimensions of each detected structure are defined by these borders. A parameter considered in our detection is the minimum number of stars included in an over-density in order to be classified as a structure, which is $N_{\text{min}} = 5$, in line with other identification techniques (e.g., Bastian et al. 2007). This criterion eliminates the detection of random stellar congregations, so-called asterisms. Additionally, we consider as real those structures which appear in at least two

7 The noun “isopleth” is used to define every line on the map that connects points having equal numeric value of surface stellar density; Origin from the Greek “isoiplēthēs”, equal in number.
Table 2. Survey of the young stellar structures in NGC 6503. Explanations on the characteristic parameters of the structures are given in the text (Sect. 3.1). Only the first 15 records of the survey are shown here for reference. The complete catalog of 244 structures is available on-line at LEGUS site https://legus.stsci.edu/.

| ID | σ | R.A. [J2000 (deg)] | Decl. | N⋆ | Radius (pc) | ξ | ρ (×10⁻³) (stars pc⁻²) | 275tot (mag) | 336tot (mag) | 438tot (mag) | Family id. | Group id. |
|----|---|-------------------|------|----|-------------|---|------------------------|-------------|-------------|-------------|-----------|-----------|
| 1  | 1 | 267.366870        | 70.139878 | 4388 | 1158.7      | 2424.7 | 2.093 | 1.04 | 13.73 | 14.21 | 15.13 | 1 | 1 |
| 2  | 1 | 267.360523        | 70.146816 | 2189 | 846.0       | 1572.5 | 1.859 | 0.97 | 14.68 | 15.13 | 16.02 | 2 | 21 |
| 3  | 1 | 267.399790        | 70.138429 | 643  | 452.4       | 670.2  | 1.481 | 1.00 | 16.26 | 16.67 | 17.57 | 3 | 36 |
| 4  | 1 | 267.324116        | 70.151839 | 206  | 255.9       | 353.4  | 1.381 | 1.01 | 17.12 | 17.30 | 17.93 | 4 | 37 |
| 5  | 1 | 267.315085        | 70.150455 | 114  | 200.6       | 375.2  | 1.271 | 0.91 | 17.51 | 18.55 | 19.43 | 5 | 38 |
| 6  | 1 | 267.360586        | 70.148180 | 97   | 194.2       | 252.0  | 1.298 | 0.83 | 17.12 | 17.53 | 18.86 | 6 | 39 |
| 7  | 1 | 267.317336        | 70.151773 | 77   | 176.4       | 238.8  | 1.354 | 0.80 | 17.69 | 18.02 | 18.94 | 7 | 40 |
| 8  | 1 | 267.325148        | 70.143415 | 62   | 148.4       | 212.8  | 1.434 | 0.91 | 17.51 | 17.92 | 18.78 | 8 | 41 |
| 9  | 2 | 267.392443        | 70.135681 | 1197 | 481.5       | 840.8  | 1.746 | 1.64 | 15.09 | 15.44 | 16.46 | 1 | 1 |
| 10 | 2 | 267.349994        | 70.143415 | 1007 | 419.9       | 1022.3 | 2.435 | 1.82 | 15.09 | 15.62 | 16.57 | 1 | 2 |
| 11 | 2 | 267.399977        | 70.138375 | 529  | 362.4       | 547.9  | 1.512 | 1.28 | 16.49 | 16.90 | 17.78 | 3 | 36 |
| 12 | 2 | 267.345413        | 70.149136 | 515  | 357.1       | 751.0  | 2.103 | 1.29 | 16.23 | 16.71 | 17.54 | 2 | 21 |
| 13 | 2 | 267.365176        | 70.140677 | 319  | 278.0       | 551.5  | 1.984 | 1.32 | 16.57 | 17.05 | 17.87 | 1 | 3 |
| 14 | 2 | 267.376002        | 70.150555 | 319  | 261.3       | 604.3  | 2.313 | 1.49 | 16.78 | 17.28 | 18.19 | 2 | 22 |
| 15 | 2 | 267.333935        | 70.151913 | 181  | 205.8       | 391.5  | 1.902 | 1.37 | 17.49 | 17.84 | 18.78 | 2 | 23 |

The catalog of identified stellar structures in NGC 6503 consists of 244 systems, revealed at significance levels between 1 and 16σ. In Column 1 we label objects with their identifier, while in Column 2 we provide the detection density threshold in σ, as measured for every object. Columns 3 and 4 give the celestial coordinates of the structures’ barycentres, which correspond to their KDE density centres. Column 5 shows the number of blue stars included within the borders of every structure, as defined by the corresponding isopleth. A measure of the size of each system is the so-called effective radius (e.g., Carpenter 2000) or equivalent radius (e.g., Román-Zúñiga et al. 2008), defined as the radius of a circle with the same area as the area covered by the system. We provide two measurements for this radius: (1) The radius determined by the area AcH enclosed by the convex hull of the system (r_eff ≡ √A_c/π, Col. 6) and (2) the radius defined by the area A_max enclosed by the largest circle that encompasses the entire system (r_max ≡ √A_max/π, Col. 7), equivalent to the half of the distance between the two furthest sources in the system.

Figure 8. Left panel. Iso-density contour plot of the surface stellar density map constructed with the KDE method with a kernel of FWHM ~ 80 pc, for the blue stars identified in both filters F275W and F336W. Right panel. Chart of the identified stellar structures, i.e., those which meet the criteria of including ≥ 5 stellar members, and appearing in at least two surface density significance levels. Systems borders are drawn with circles representing their sizes, r_eff, as determined at the various significance levels, where they appear. In both panels different colors are used for the isopleths and apertures of systems found at different significance levels in steps of 1σ.
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The ratio of these two radii provides a characterization of the elongation of each system $\xi \equiv r_{\text{max}}/r_{\text{eff}}$, which we provide in column 8. Schmeja & Klessen (2006) showed that the estimation of the system’s area from its convex hull makes $\xi$ a reliable measure for the elongation, since it excludes the possibility that fractal substructure in an otherwise spherical structure can lead to unrealistic elongations. These authors also determined the increase of $\xi$ with increasing axis ratios of elliptical distributions; A circular distribution (axis ratio = 1) has $\xi \simeq 1$. From the number of blue stars enclosed within the borders of each structure (Col. 5) and its radius $r_{\text{eff}}$ (Col. 6) we also have a measurement of the surface stellar density of each structure, which is given in column 9 of Table 2. These stellar structures contain blue stars in the whole observed magnitude range of $F336W \lesssim 26$. The total brightness of each system, calculated from its stellar population in the three blue filters (i.e., F275W, F336W and F438W) is given in columns 10, 11 and 12 respectively.

A histogram of the numbers of all detected structures per detection density level up to $10 \sigma$ is shown in Fig. 9. This distribution peaks at the detection levels of 2, 3 and $4 \sigma$, which determine the vast majority of the detected structures. Higher detection levels mostly reveal the high-density peaks of the same structures, as well as their “offspring” structures which they break into. A chart of the detected structures is shown in Fig. 8 (right panel). In this map the structures are indicated by circles equivalent to their measured sizes. Apertures with radii equal to $r_{\text{eff}}$ of the detected systems are drawn in different colors according to the density threshold where the structures were detected (as in the isopleth plot at the left panel of the figure).

As shown in Table 2, and can be seen in the maps of Fig. 8, there are eight large structures identified at the $1 \sigma$ level. Two of them extend to the $2 \sigma$ level, and appear isolated, i.e., not related to other large structures. Three additional $1 \sigma$ structures appear also in few intermediate density levels ($\geq 2 \sigma$), two of them also showing a second sub-structure. The remaining three $1 \sigma$ structures are extended stellar concentrations that qualify as large stellar super-complexes with dimensions between 1 and 2 kpc. These structures comprise the vast majority of identified stellar structures, which are themselves multiple concentrations seen at higher density levels. The $2 \sigma$ concentrations most probably correspond to the so-called stellar aggregates and stellar complexes of the galaxy, while the structures detected at the $3 \sigma$ level are actually members of these concentrations, fulfilling the typical image of hierarchical structuring of stars on a galaxy-scale. Even higher density detections ($\geq 4 \sigma$) correspond to the more condensed stellar density peaks of the larger structures.

We define each of the structures revealed at the $1 \sigma$ level as a family of structures. We also determine as members of a group structures that are “spawned” from the same single structure. In total we identify 82 groups divided into eight families, three of which are large super-complexes that encompass 95% of all revealed stellar structures. The family and group numbers of each structure are given inCols. 13 and 14 of Table 2 respectively.

3.2 Structure Tree

In this section we visualize the “family ties” of high density structures in relation to their parental super-structures. The “breaking” of super-structures of specific low density level into several, smaller and denser sub-structures at higher density levels, provides the first evidence of a morphological hierarchy in the way young stars are clustered across the ring. An intuitive way to illustrate hierarchical structures is through the so-called dendrograms, introduced as “structure trees” for the analysis of molecular cloud structure by Houllahan & Scalo (1992), refined by Rosolowsky et al. (2008). A dendrogram is constructed by identifying in the stellar surface density map connected structures found at different density thresholds, while keeping track of the connection to “parent structures” (on a lower level) and “child structures” (on the next higher level, lying...
Table 3. Demographics of the young stellar structures revealed in the KDE surface stellar density map at various significance levels. In Col. 1 the detection levels (in \( \sigma \)) are given. The parameters shown are the average size (Col. 2) and size dispersion (Col. 3), the average elongation (Col. 4), and the average surface stellar density (Col. 5) of the structures in each density level. The corresponding total stellar numbers (\( N_\star \), Col. 6) and UV magnitudes (\( m_{275} \), Col. 8) per density level are also given, along with the corresponding fractions of these parameters over the total observed number of stars, \( N_\star/\sigma \), and total observed stellar UV flux, \( m_{275, tot} \), shown in Cols. 7 and 9, respectively. There are no dispersions given for the parameters in the last three entrances, because there is only one object found in each of the corresponding density levels.

| Detection level (\( \sigma \)) | Size \( S \) (pc) | Size dispersion \( \sigma_S \) | Elongation \( \xi \) | \( \bar{\rho} \times 10^{-3} \) (stars pc\(^{-2} \)) | \( N_\star \) | \( f_\star \) | \( m_{275} \) (mag) | \( m_{275, tot} \) |
|-----------------------------|----------------|-------------------|--------|-------------------|--------|---------|-------------|-------------|
| 1                           | 858            | 751               | 1.60 ± 0.30 | 0.93 ± 0.08 | 7776   | 0.606   | 13.2       | 0.699       |
| 2                           | 295            | 210               | 1.49 ± 0.39 | 1.46 ± 0.20 | 6126   | 0.477   | 13.3       | 0.603       |
| 3                           | 202            | 132               | 1.32 ± 0.29 | 1.97 ± 0.26 | 4371   | 0.341   | 13.6       | 0.474       |
| 4                           | 148            | 104               | 1.30 ± 0.33 | 2.94 ± 1.00 | 3070   | 0.239   | 13.9       | 0.367       |
| 5                           | 127            | 94                | 1.17 ± 0.20 | 3.46 ± 0.81 | 2043   | 0.159   | 14.2       | 0.283       |
| 6                           | 129            | 97                | 1.22 ± 0.28 | 3.86 ± 1.04 | 1400   | 0.109   | 14.5       | 0.211       |
| 7                           | 121            | 88                | 1.18 ± 0.26 | 4.80 ± 1.54 | 907    | 0.071   | 14.8       | 0.151       |
| 8                           | 142            | 92                | 1.21 ± 0.29 | 5.20 ± 2.11 | 623    | 0.049   | 15.1       | 0.117       |
| 9                           | 179            | 45                | 1.34 ± 0.19 | 4.55 ± 0.52 | 475    | 0.037   | 15.3       | 0.099       |
| 10                          | 146            | 47                | 1.31 ± 0.23 | 5.15 ± 0.90 | 362    | 0.028   | 15.5       | 0.082       |
| 11                          | 125            | 63                | 1.32 ± 0.39 | 6.08 ± 0.80 | 243    | 0.019   | 15.8       | 0.063       |
| 12                          | 138            | 6                 | 1.47 ± 0.48 | 6.05 ± 0.60 | 179    | 0.014   | 16.0       | 0.052       |
| 13                          | 110            | 10                | 1.57 ± 0.56 | 6.43 ± 0.56 | 121    | 0.009   | 16.2       | 0.044       |
| 14                          | 99             | -                 | 1.14       | 7.50         | 57     | 0.004   | 16.8       | 0.024       |
| 15                          | 81             | -                 | 1.24       | 8.99         | 45     | 0.004   | 17.0       | 0.021       |
| 16                          | 59             | -                 | 1.07       | 6.87         | 18     | 0.001   | 17.4       | 0.014       |

within the boundaries of its parent). In the dendrogram of a geometrically perfect hierarchy each parent would branch out into the same number of children at each level.

We construct the dendrogram of the stellar structures detected in NGC 6503 at various density thresholds, up to the highest level of 16\( \sigma \) above the background density. In this dendrogram, shown in Fig. 10, the structures found at each density level are represented not only by the ‘leaves’ that end at the particular level, but by all branches present at that level. For example, at the 3\( \sigma \) level, there are 50 branches of the dendrogram (regardless whether they end at this level, continue to a higher level as a single system, or split into two or more branches), corresponding to the 50 detected stellar structures. This dendrogram demonstrates that most structures split up into several sub-structures over few levels. The combination of this dendrogram with the maps of Fig. 8 illustrates graphically the hierarchical spatial distribution of young stars in NGC 6503.

### 3.3 Parameters demographics

Our survey of young stellar structures (Table 2) covers a variety of systems, starting with those including the minimum of 6 to 8 stellar members with dimensions \( \lesssim 40 \) pc (e.g., structures 148 and 226), up to those with maximum size of more than 1 kpc, including over 1000 stars (i.e., structure 9). The parameters demographics of the stellar structures, revealed at various significance levels, are given in Table 3. These parameters include the average size (Col. 2) and size dispersion (Col. 3), the average elongation (Col. 4), and the average surface stellar density (Col. 5) of all structures found in each density level. The total numbers of stars and total UV magnitudes of all structures in each level are also given (Cols. 6 and 8 respectively). Table 3 provides the fraction of stars at each significance level relative to the total young stellar sample (Col. 7), and the corresponding stellar UV flux fraction relative to the total observed UV flux per detection level (Col. 9).

In general, almost all parameters given in Table 3 show a dependence on the detection level\(^8\). The average size (and its dispersion), decreases with increasing density level, with a plateau and a small bump toward relatively larger sizes for levels between 6 and 12\( \sigma \). The mean and median of the sizes reported in Col. 2 of Table 3 equal to 185 pc and 138 pc respectively. Both the total stellar number and total UV brightness show a systematic correlation with the detection significance level, with larger and sparser stellar structures hosting higher stellar numbers and UV brightness. This agreement in the trends of these two parameters can be directly explained by the almost one-to-one correlation between their values, as can be derived from the data of Table 3.

A systematic dependence on level also exists for the fraction \( f_\star \) of stars included in every density level over the total observed number of blue stars. This fraction changes from 60% within the 1\( \sigma \) structures to \( \sim 1\% \) at the highest density level (corresponding to one structure identified at 16\( \sigma \) level). However, while the same dependence on significance level is found for the fraction of UV emission per level relative to that observed from all blue stars (Col. 9 in Table 3), this trend is steeper than that for stellar fraction. In particular, while \( \sim 60\% \) of the stars (i.e., the 1\( \sigma \) structures) produces \( \sim 70\% \) of UV, this correlation changes for structures with higher densities. For example, at 7\( \sigma \) significance, where \( \sim 7\% \) of the total young population resides, \( \sim 15\% \) of the total UV brightness is produced. At even higher levels only a few \( \% \) of the total young stellar population in NGC 6503 emits few \% of the total UV light (see Table 3). This comparison suggests that compact structures, identified at higher density levels, encompass on average the UV-brightest stars in the galaxy. This result is elaborated more in Sect. 4.3, where we discuss the clustering behavior of blue stars as a function of their brightness.

The size distribution of all detected structures constructed by

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\(^8\) Among all parameters, only elongation, \( \xi \), seems to be independent of the level of detection for the structures.
binning them according to the logarithm of their dimensions is shown in Fig. 11. Dimensions are given in physical units (pc), and they are derived from the effective radii of the structures. A functional fit of this histogram to a normal distribution, drawn with a red line, shows that the dimensions of the detected systems are clustered around an average of \( \sim 120 \text{ pc} \) with a standard deviation of \( \sim 40 \text{ pc} \). This is an interesting result, considering that this size is comparable (but still to the upper limit) to the typical size for giant molecular clouds (e.g., Cox 2000; Tielens 2005). It has been also pointed out earlier by various authors (e.g., Efremov et al. 1987; Ivanov 1996) that young stellar associations in different galaxies have dimensions that average at \( \sim 80 \text{ pc} \), comparable to the 40% percentile of our distribution. Considering that the determination of structures dimensions is sensitive to the resolution of their detection, the question if this length-scale would represent a characteristic scale for star formation is still open (e.g., Bastian et al. 2007; Gouliermis 2011). Moreover, the detection resolution seems also to determine the peak of this distribution. By repeating the analysis after smoothing by smaller scales (between 70 and 40 pc), we found that indeed the peak moves at smaller values (from \( \sim 120 \text{ pc} \) at 80 pc resolution to \( \sim 70 \text{ pc} \) at 40 pc resolution).

From this fit it is seen that the size distribution of our structures is not lognormal. There is a deficiency in the number of small-scale structures at the left hand tail of the distribution, which falls below the Gaussian. This deficiency is most probably due to incompleteness introduced by the smoothing, since it seems to be more important for the distributions derived with larger smoothing kernels. In addition, the right hand part of the distribution, for structures with sizes larger than \( \sim 120 \text{ pc} \), is more extended than the Gaussian and has a power-law shape. This is demonstrated with a power-law fit, shown with the blue dashed line, which indicates a scale-free behavior. It is interesting to note that the power-law behavior of the right part of the distribution seems to be independent of the smoothing kernel, having almost the same exponent of \( \sim -1.5 \) for distributions produced after smoothing with smaller kernels.

This power-law behavior is further supported by the cumulative radius distribution of the star-forming structures found in different density thresholds (levels), shown in Fig. 12. The distributions of structures found only within five indicative detection significance levels are shown to avoid confusion. The cumulative distributions functions of radii for all structures or those including low density structures (blue symbols) show power-law tails at large radii. The distributions for samples constrained to high-density structures (green and yellow symbols) have shapes closer to lognormal.

### Figure 11.
Histogram of the size distribution of all detected young stellar structures in NGC 6503, with a logarithmic bin size of 0.075. Sizes are defined as \( 2 \times r_{\text{eff}} \) and are given in pc. The distribution peaks at a size of \( \sim 130 \text{ pc} \). The best-fit Gaussian (red line) peaks at an average size of \( \sim 120 \text{ pc} \), close to the smoothing radius used for the structures identification of 80 pc (indicated with the vertical grey dashed line). This fit shows that the right hand part of the size distribution, corresponding to the resolved structures, behaves almost like a power law (indicated by the blue dashed line).

### Figure 12.
Radius cumulative distribution functions derived from the NGC 6503 survey of star-forming structures. The results for five different detection density levels are shown. The distribution of structures found at levels \( \geq 1 \sigma \) (purple symbols) corresponds to the total sample. This distribution and those for samples including low-density structures (blue symbols) show power-law tails at large radii. The distributions for samples constrained to high-density structures (green and yellow symbols) have shapes closer to lognormal.

#### 3.4 Parameters correlations

The correlation of physical parameters of star-forming structures provide insight to the processes and conditions that dictate their formation and evolution. If these systems do trace gas structure, then their fundamental properties should be tightly related to those of GMCs, and should have comparable correlations, which nevertheless should also be affected by the galactic dynamics and the transformation of gas to stars. In this section we explore the correlation of the structural parameters measured for the discovered stellar structures in order to identify the factors that determined those relations. The basic parameters considered are those directly measured from our data, namely the size of the structures (determined as two times their \( r_{\text{eff}} \)), their surface stellar density, \( \mu \) (in stars pc\(^{-2}\)), the number of member stars, \( N_* \), and their total UV brightness, \( m_{275} \), in magnitudes.

We also make a rough estimate of the volume stellar density of each structure, \( \rho \) (in stars pc\(^{-3}\)), assuming spherical symmetry, as \( \rho \propto \mu^2 r_{\text{eff}}^2 \). This estimation for elongated structures, which are typically the larger low-density structures, is expected to suffer from the lack of symmetry. However, we determine the volume density within the effective radius and not the maximum radius, which defines the largest aperture covered by the structure, and includes a large fraction of empty areas. Thus we eliminate to a large extent the bias in the volume density estimation due to asymmetries.
3.4.1 Stellar Number Correlation with Radius

To measure molecular clouds structure Kauffmann et al. (2010a,b) established the correlation between mass and radius for clouds in the solar neighborhood. Using the survey of stellar complexes of the present work, we can test such a relation for large star-forming structures. As a first-order approximation, the number of stars in a group is directly proportional to its mass, assuming all groups are of a similar age and as long as sampling effects of the mass function are not significant. In Fig. 13 we show the measured radii, $r_{\text{eff}}$, versus the number of stars within each identified structure. The solid line in the plot represents the least-square linear fit to the logarithms of the data. While there is no significant scatter in the whole sample, the relation for systems identified at lower density thresholds is different than that for those found at higher densities with indexes varying between about 1.5 and 2. This trend is consistent to that found by Gouliermis et al. (2003) for 494 stellar associations and open clusters in the Large Magellanic Cloud.

A single power-law mass-size relation with a fractional index is expected from a fractal distribution of young stars (in agreement with, e.g., Elmegreen & Falgarone 1996). A uniform two-dimensional distribution of stars would correspond to a constant surface stellar density, and therefore the number of stars (or mass) would be proportional to the radius (or size) as $\propto \pi r^2$. Thus, a uniform stellar distribution would also produce a power-law behavior but with an index 2 due to the dependence of surface to radius. This relation is shown with a dashed line in Fig. 13. Concerning structures in our sample, the correlation between number of stars and radius for the low-density structures (up to $3\sigma$) has slopes very close to 2, indicating that these systems have constant surface densities. On the other hand, the high-density systems show correlations different to uniform distribution with slopes between 1.5 and 1.8 (all slopes have about 1% uncertainty), consistent with self-similar stellar distributions. Such a dependence of the mass-radius relation to the detection threshold is also found in M 33 by Bastian et al. (2007, who used a different technique), summarizing the differences in the trends of properties for different types of stellar systems. We explore these differences in the following section.

3.4.2 Surface and Volume Stellar Density Correlations with Size

Correlations between the surface stellar density $\rho$ and volume stellar density $\rho$ of the young stellar structures to their sizes $L$ (in pc) are shown in Fig. 14. Points are colored according to the detection surface density level (in $\sigma$) of the corresponding structures. In the top panel it is shown that the surface density of low-density (large) structures, in particular those detected at levels lower than $\sim 5\sigma$, is almost constant, in contrast to the high-density systems, where a correlation of surface density and size is more apparent. On the other hand, volume density, plotted in the bottom panel, shows a clear correlation with size for all structures, with different exponents for structures found at different density levels. The shaded area represents Larson’s 3rd relation, which correlates the sizes of molecular clouds to their volume densities. The correlation found for our stellar structures seems to comply with this relation, although it has a trend to be on average steeper (fit is shown with the thick line). The correlation for systems detected at levels lower than $\sim 5\sigma$ is more compatible to Larson’s relation than those for systems found at levels $\gtrsim 5\sigma$, which are steeper (see also Fig. 15).

Figure 13. Correlation of number of stars with the radius, $r_{\text{eff}}$, of the detected stellar structures. The different colored points represent structures identified at different density levels, indicated by the legend. The black solid line is the best fit through all the data, with a slope of 1.6. This fit for the various significance structures varies in agreement with previous investigations. This power-law behavior is expected from a fractal stellar distributions. A slope of 2, corresponding to constant stellar surface density, is shown with the grey dashed line to indicate how a uniform distribution appears in this correlation.

Figure 14. Surface stellar density (top) and volume stellar density (bottom) correlation with the size of the detected stellar structures. In the top panel it is shown that the surface density of low-density (large) structures, in particular those detected at levels lower than $\sim 5\sigma$, is almost constant, in contrast to the high-density systems, where a correlation of surface density and size is more apparent. On the other hand, volume density, plotted in the bottom panel, shows a clear correlation with size for all structures, with different exponents for structures found at different density levels. The shaded area represents Larson’s 3rd relation, which correlates the sizes of molecular clouds to their volume densities. The correlation found for our stellar structures seems to comply with this relation, although it has a trend to be on average steeper (fit is shown with the thick line). The correlation for systems detected at levels lower than $\sim 5\sigma$ is more compatible to Larson’s relation than those for systems found at levels $\gtrsim 5\sigma$, which are steeper (see also Fig. 15).
Hierarchical Star Formation across the ring galaxy NGC 6503

Figure 15. Relation between the index $\alpha$ in the relation $\rho \propto L^{-\alpha}$ for the detected stellar structures and the corresponding detection density level (in $\sigma$). This plot shows a dependency of the correlation between volume stellar density and size of the structures to their detection limits. The size of large, low-density structures (found at levels up to $\sim 5\sigma$), seems to be correlated to their corresponding volume density in a fashion comparable to or flatter than Larson’s third relation (indicated by the horizontal dashed dotted line), while high-density, smaller structures have clearly steeper correlation between their sizes and volume densities.

The volume density-size correlation, shown at the bottom panel of Fig. 14, can be represented for all points by a power law of the form $\rho \propto L^{-\alpha}$ with index $\alpha = -1.33$, established with a linear fit to the data (grey thick line in the plot). However, as can be seen in the plot, this slope differs for structures identified at different density detection levels, with low-density (and larger sizes) structures following a more shallow relation than the high-density structures. We demonstrate the dependency of the exponent $\alpha$ for our stellar structures to their detection density level in Fig. 15, where we plot $\alpha$, as found for every group of structures, as a function of the corresponding density threshold (in $\sigma$).

The correlation of volume densities and sizes we establish for the detected stellar structures can be directly compared to that determined by Larson (1981) for molecular clouds, and confirmed by other authors (e.g., Myers 1983; Solomon et al. 1987; Heyer et al. 2009). The so-called Larson’s third relation exhibits the tendency of the mean cloud volume density, $\rho$ (in cm$^{-3}$) to scale inversely with the cloud size, $\rho \propto L^{-1.1}$. Larson’s exponent of $-1.1$ is indicated by a horizontal dashed-dotted line in Fig. 15. This exponent coincides with $\alpha$ determined in our correlation for structures at density levels between $\sim 5\sigma$ and $6\sigma$.

Larson’s relation implies that clouds have approximately constant column densities, i.e., the same significance levels relative to a fixed background, reflecting the identification of molecular clouds as constant-density objects (Solomon et al. 1987). In the case of our stellar structures, as seen in Fig. 14 (top panel) this does not apply, since the identification of these structures is made at various density levels. Only structures detected at low density levels ($< 5\sigma$) have surface stellar densities almost independent of their size, and $\rho(L)$ relations comparable to Larson’s relation ($\alpha$ values below the 1.1 limit in Fig. 15). In contrast, high-density structures have steeper $\rho(L)$ relations, because their $\mu(L)$ relations are not flat. These steeper correlations are the result of the higher densities of structures that tend to be smaller, and reflect the steeper density profiles of these systems (as $\propto r^{-\alpha}$ for spherical systems) in relation to those of low density.

Figure 16. (a) Cumulative radial stellar density profile of the blue stellar population in NGC 6503. This profile exhibits the ring-shaped distribution of the youngest populations, highlighting the very low abundance of young stars in the inner parts of the galaxy. In this profile the star-forming ring peaks at a radial distance $\sim 72'' +10'\!-\!8'\!-\!10''$ (\sim $1.85\!+\!0.26'\!-\!0.21$ kpc). Exponential disk fits of the profiles are shown with dashed lines. These fits, equivalent to those on the surface brightness profiles previously presented (Bottema 1989; Freeland et al. 2010), do not fit the outskirts of the observed distributions, which are fitted with different exponential profiles (shown with dotted lines).

(b) The corresponding profile of the red population shows a different behavior, with a much lower deficiency of stars in the inner parts, and a peak at somewhat larger radial distance of $\sim 81'' +11'\!-\!10'\!-\!10''$ (\sim $2.09\!+\!0.28 -0.26$ kpc). Exponential disk fits of the profiles are shown with dashed lines. These fits, equivalent to those on the surface brightness profiles previously presented (Bottema 1989; Freeland et al. 2010), do not fit the outskirts of the observed distributions, which are fitted with different exponential profiles (shown with dotted lines).

It should be noted, though, that while our structures for detection levels $\leq 6\sigma$ have sizes larger than $\sim 100$ pc, most of the molecular clouds in the Milky Way have sizes of that order or less.
4 THE GLOBAL CLUSTERING OF YOUNG STARS IN NGC 6503

4.1 Stellar Surface Density Profiles

We construct the radial stellar surface density profile across the observed field-of-view of the blue stellar population of NGC 6503, by counting the stars in concentric annuli around the rotational center of the galaxy (defined in Sect. 2.2). The annuli are determined in logarithmically increasing radial distances, so that the width of each annulus increases with its radial distance. The cumulative radial stellar surface density profile of the blue stellar population of NGC 6503 is shown in Fig. 16 (top panel). In this figure the profile of the red population, as this stellar sample is determined in Sect. 2.1, is also shown for comparison (bottom panel).

From these profiles it is seen that blue and red stars follow quite different radial distributions around the center of the galaxy. Specifically, the blue stars show a very strong peak at the center, while the LINER is located. The surface density drops rapidly within a very short radial distance (∼ 14′′ ~ 0.35 kpc) and then rises sharply, reaching a peak at a distance of ∼ 1.9 kpc (Fig. 16, top). The distribution of the red stars shows also central densities lower than those at larger radii, but their deficiency in the inner part of the galaxy is not as severe as for the blue stars. Nevertheless, the surface density of red stars rises slowly outwards and peaks at a distance of ∼ 2.1 kpc, a distance comparable to that where the blue stellar distribution peaks. It drops again at larger distances, but far more slowly than that for the blue stars (Fig. 16).

We evaluated the deficiency of stars in the galaxy center in terms of incompleteness in our photometry due to crowding. Our artificial star experiments showed that the lack of blue stars in the circumnuclear region is not due to photometric incompleteness. The number of UV sources is apparently very small and thus the blue stellar sample is flux- rather than crowded-limited even at the center of NGC 6503. The situation is different in the optical filters, as for example in the F555W filter, which is found to be clearly more incomplete in the center than in the periphery. As a consequence, the circumnuclear decreases seen in the stellar profiles of Fig. 16 are probably due to different reasons. While the central lack of blue sources is due to the real absence of a significant (and dense) young population in the center, the central deficiency of red stars is due (at least in part) to incompleteness.

Both surface density profiles shown in Fig. 16, and in particular that of the blue young population, track the inner star-forming ring of the galaxy with a peak at ∼ 2 kpc. This is in line with the previously published radial luminosity profile of NGC 6503, which is proposed by Bottema (1989) as a typical example of Type II (Freeman 1970), i.e., of the form I(R) < I_0 e^{-αR} for a radial distance interval not far from its center. At larger radial distances both our blue and red profiles drop in an exponential fashion, as shown by the fitted exponential disk models plotted with dashed lines on the observed profiles of Fig. 16. These exponential disk models agree with those describing the smooth surface brightness profile of the galaxy in various wavelengths for radial distances up to about 3 kpc (Freeland et al. 2010). However, in order to describe the distribution of stars at the outskirts of our stellar density profiles for radial distances ∼ 4-5 kpc, fits to additional exponential disks were considered. These fits are plotted in the profiles of Fig. 16 with dotted lines.

A Type II profile produced by a stellar deficit in the inner parts of the galaxy may be caused by (1) dust extinction, (2) an inner truncated disk, or (3) a ring of bright stars. An exponential disk and a truncated light profile due to dust extinction are proposed for NGC 6503 by Bottema & Gerritsen (1997), but a statistical study on several galaxies showed that the possibility of dust extinction causing Type II profiles is inconclusive (MacArthur et al. 2003). Concerning the other two possibilities, a truncated disk in NGC 6503 implies that there is a genuine mass deficit in the inner region of the galaxy, and a stellar ring relates to bar formation. The dynamics of NGC 6503 was modeled by Puglirelli et al. (2010) with a Bayesian inference by assuming a Kormendy (1977) truncated disk. These authors found that the disk of NGC 6503 is indeed equally well-fitted by either mechanisms, i.e., an inner-truncated profile, or a ring formation by a bar.

The inner drop of stellar density seen in our profiles also agrees with a sharp decrease in the velocity dispersion (σ-drop) for radii smaller that ∼ 10′′ (∼ 0.25 kpc) observed by Bottema (1989) with a minimum at a distance remarkably similar to that of our blue stellar surface density profile10. Modeling of the dynamics of the galaxy, however, could not provide a reliable theoretical description of the inner drop in dispersion (Bottema & Gerritsen 1997). In their study Puglirelli et al. (2010) found that the mass-luminosity ratio of NGC 6503 pseudobulge is lower than that in the disk, suggesting the presence of a dynamically-cold star-forming component that is probably responsible for the velocity dispersion drop, in agreement with previous studies (Wozniak et al. 2003; Comerón et al. 2008). The existence of an inner nuclear bar could also produce the σ-drop, as has been reported in the literature for some cases (de Lorenzo-Cáceres et al. 2008). Indeed, while the ring of NGC 6503 is not considered to have the classical aspect of dynamically induced resonance rings (Mazzuca et al. 2008), it is likely an inner ring and not a nuclear ring, possibly caused by a strong end-on bar which is embedded inside it (Knapen et al. 2006; Freeland et al. 2010).

4.2 Hierarchical Structure in Stellar Clustering

As discussed in Sect. 3, young stellar structures in NGC 6503 are assembled in a hierarchical fashion, in the sense that there are condensed structures belonging to larger looser ones, which themselves are part of even larger low-density stellar systems. This behavior, demonstrated visually in the dendrogram of Fig. 10, suggests that the morphology of young stellar cluster assembling across the whole galaxy is self-similar. In this section we quantify the young stellar clustering behavior with the use of the two-point correlation function and we determine the time-scale within which stellar clustering sustains its behavior.

The spatial distribution of stars can be quantified with the construction of the two-point correlation or autocorrelation function (ACF), which is a measure of the degree of clustering in the spatial distribution, ξ(r), of a sample of sources (Baugh & Murdin 2000). This method, introduced by Peebles (1980) in cosmology, has been successfully used for characterizing the stellar clustering behavior in star-forming regions in the Milky Way (e.g. Gomez et al. 1993; Larson 1995), as well as that of stellar populations and star clusters in remote galaxies (e.g., Bastian et al. 2009; Scheepmaker et al. 2009). In this study we apply the method to the resolved stellar population of a whole galaxy, following the prescription by Gouliermis, Hony & Klessen (2014). The innovation of our treatment lies on the use of the de-projected positions of the stars, eliminating the effect of projection on the measures of stellar pair-separations.

10 The velocity dispersion drop in NGC 6503 is also discussed by Bottema & Gerritsen (1997) and Puglirelli et al. (2010)
to a power spectra analysis. Both statistics are used in studies of cosmological large-scale structure (e.g., Szapudi et al. 2005) and of density structure in the turbulent interstellar matter (e.g., Federrath et al. 2009).

In Fig. 17 we show the ACF of the blue stellar sample observed in NGC 6503 (blue line). For comparison we also show in red the ACF of the red stellar sample. Both functions were possible to be constructed for separations down to 0.75 seconds of arc, corresponding to physical scales of $\sim 20$ pc. For smaller scales, there is an “anti-correlation” of the ACF with separation, which drops rapidly towards the smallest separations. Apparently, photometric confusion dominates at scales smaller than 20 pc, a limit which, thus, sets the resolution of our analysis with these data. The ACF, being determined through the pair-separations between all stars, can be calculated within a specific maximum length-scale limit, determined by the length of the unavoidably finite observed field-of-view. In practice, the continuously larger stellar pair-separations calculated for each star in the sample eventually reach the edge of the observed field, beyond which there are no stars available for the calculation. If the ACF determination is not corrected for this edge-effect, the calculation will be incorrect at large scales with its value dropping to the unrealistic $1 + \xi(r) < 1$.

We correct our calculations by masking the pair-separations measurements within the borders of the observed field-of-view, as prescribed by Gouliermis, Hony & Klessen (2014). This treatment corrects the stellar densities of Eq. (2) by using the true area surface constrained for the larger length-scales by the masking. This correction allows the ACF to fall smoothly to the limiting value of 1 at scales comparable to the edge of the observed field as shown in the plots Fig. 17. In effect, the sharp drop to the limiting value of 1 sets the maximum length-scale within which the ACF is reliable. We denote in Fig. 17 with the greyed area the remaining unreliable length-scale range, which we do not take into account in our further analysis. It is interesting to note that the maximum scale, where a trustworthy ACF is calculated is $\approx 100''$ ($\approx 2.6$ kpc), far shorter than the complete extent of the de-projected field-of-view.

We determine the exponent $\eta$ of the ACF through a linear regression on the log-log plots of Fig. 17 by applying a Levenberg-Marquardt nonlinear least square minimization fit (Levenberg 1944; Marquardt 1963). In this figure it is seen that both the blue and red stellar samples in NGC 6503 have a single power-law stellar separations dependency of their ACF. The exponent $\eta$, however, is quite different between the two samples. The ACF of the young blue sample has a slope $\eta \approx -0.30$, corresponding to a fractal dimension $D_2 = 1.7$. On the other hand, the ACF of the old red stellar sample is almost flat with $\eta \approx -0.05$, i.e., $D_2 = 1.95$ (both slopes found with very small fitting errors).

There are two conclusions, connected to each other, one can derive from these results. (1) The steep monotonic ACF of the blue stars, corresponding to a fractal dimension much smaller than the geometrical dimension of 2, suggests a fractal, i.e., self-similar distribution of young stars across NGC 6503. On the other hand, the flatter ACF for the red stars, with a fractal dimension very close to the geometric dimension, clearly implies a distribution for the old population, which is very well spread, almost equivalent to a uniform (random) spatial distribution. (2) The absolute values for

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**Figure 17.** The two-point correlation (autocorrelation) function constructed for the two stellar samples of interest, revealed with LEGUS photometry in NGC 6503. The ACF of the young blue stellar population detected in F275W and F336W filters is shown with the blue line, while that of the older red stars selected in filters F555W and F814W is drawn with the red line. The grey flat dotted line represents the autocorrelation function of a uniformly distributed stellar population. From this plot it is shown that the young blue population follows a clustering behavior, which is quite different than that of the evolved red stars, being more clustered than the latter. The ACF of the red stars, being almost flat, demonstrates that these stars are quite more distributed than the blue. The shaded area designates the length-scale, where the ACF calculation is unreliable – and thus not considered in our analysis – due to the edge-effect introduced by the limited field-of-view.

The constructed ACF is, thus, the true two-dimensional ACF of the galaxy, as if it was observed face-on. We determine the ACF of the stars from their de-projected coordinates, as:

$$1 + \xi(r) = \frac{1}{n_1 N} \sum_{i=1}^{N} n_i(r), \quad (2)$$

where $N$ is the total number of stars, $n_i(r)$ is the number density of stars found in an aperture of radius $r$ centered on star $i$, and $\bar{n}$ is the average stellar number density. Uncertainties of this function are given by:

$$\delta(1 + \xi(r)) = \sqrt{N} \left( \frac{1}{2} \sum_{i=1}^{N} n_p(r) \right)^{-1/2}, \quad (3)$$

where $n_p(r)$ is the number of pairs with the central star $i$ of the current aperture, and the factor 1/2 accounts for not counting every pair twice. The function $1 + \xi(r)$, is defined so that $\bar{n}[1 + \xi(r)]d^2 r$ is the probability of finding a neighboring star in an area of radius $r$ from a random star in the sample. Therefore, for a random stellar distribution $1 + \xi(r) = 1$, while a truly clustered sample should have $1 + \xi(r) > 1$.

In a two-dimensional self-similar distribution the total number of stars $N$ within an aperture of radius $r$ increases as $N \propto r^{D_2}$, where $D_2$ is the fractal dimension of the distribution (Mandelbrot 1983). Fractal distributions have a power-law dependency of the ACF with radius of the form $1 + \xi(r) \propto r^{\eta}$ (e.g., Larson 1995), which yields from Eq. (2) $N \propto r^{\eta} \cdot r^2 = r^{\eta+2}$. The exponent $\eta$ is thus related to the fractal dimension as $D_2 = \eta + 2$. For the derivation of the three-dimensional fractal dimension $D_3$ of a distribution, dedicated simulations are shown to be required (Gouliermis, Hony & Klessen 2014). The ACF is qualitatively identical method
the ACF of the young stars, which are found systematically higher than those for the old stars at the same separation length, indicate the more “clumpy” clustering behavior of the former sample in comparison to the latter, with the young stars being systematically more clustered than the old.\(^\text{13}\)

The steep ACF of the young stellar population in NGC 6503 demonstrates that these stars are hierarchically distributed across the whole measurable extent of the galaxy, up to scales $\sim 2.5$ kpc. This implies that the hierarchical stellar structures identified in Sect. 2.3, are themselves part of an extended hierarchical distribution. Power-laws similar to that found here for the ACF of young stars in NGC 6503 are derived from power spectra of interstellar gas over a large range of environments (e.g. Elmegreen & Elmegreen 2001), demonstrating the hierarchical morphology in the gas structure with a typical fractal dimension of $D_2 \simeq 1.5$ (Elmegreen et al. 2006).

4.3 Time Evolution of Stellar Structural Morphology

The sample of blue young stellar population of NGC 6503 includes stars of various ages, and therefore represents different star formation events within the recent star formation history of the galaxy. This is demonstrated in Fig. 18, where the CMD of the blue stars is shown with indicative stellar evolutionary models for solar metallicity from the Padova grid of models (Chen et al., in prep.; see also Marigo et al. 2008; Girardi et al. 2010; Bressan et al. 2012). From the oldest isochrone it can be derived that the stellar age limit covered by our blue photometric catalog is $\gtrsim 100$ Myr (red line in Fig. 18).

The stellar age range covered in our sample allows us to address the question whether star formation is a purely clustered process, or stars are also frequently born in a distributed fashion over galactic scales by investigating the clustering behavior of young stars at sequential evolutionary stages. The ACF treatment described in the previous section for the characterization of the clustering behavior of the whole blue sample can be also applied to stars of different ages. The construction of the ACF for subsamples of different ages would thus provide an assessment of how the clustering behavior of young populations evolves with time across the whole galaxy, over the last $\sim 100$ Myr, as investigated for example in the Magellanic Clouds (Gieles et al. 2008; Bastian et al. 2009).

The derivation of ages of blue bright stars from their photometry alone is a complicated process and implies significant uncertainties for the derived parameters, even with the use of multi-band photometric measurements. In order, thus, to retain most of the original information derived from our catalogs, and to avoid introducing model-dependent uncertainties, we group the sample of blue stars in ranges defined by their magnitudes in the F336W filter. We use this selection as a proxy for dividing the stars in groups of different evolutionary stages. In CMDs as this of Fig. 18, both the turn-off (TO) and the He burning, starting immediately after the TO, correspond for every isochrone to a specific F336W magnitude which for older stars is systematically fainter. As a consequence, following the analysis by, e.g., Bastian et al. (2009), we group the stars within continuously fainter magnitude bins in order to roughly divide them into age bins. It should be noted though, that while the brightest magnitude range is populated only by the youngest stars, fainter ranges should also include stars younger than the corresponding age limit. We assess, however, that the fraction of younger contaminants is small enough not to significantly affect the canonical age of each magnitude bin.\(^\text{14}\)

4.3.1 Two-Point Correlation Functions

The blue stellar catalog is divided into eight magnitude bins, all containing equal numbers of $\simeq 1,600$ stars. Dividing the catalog into ranges of equal numbers satisfies equivalent statistical significance among all sub-samples. The limiting magnitudes of each sub-sample are shown in the first column of Table 4. The corresponding limiting age and stellar mass for each sub-sample, defined by the TO of the corresponding isochrone, are given in the second and third columns of the table respectively. We construct the ACF for each of the blue sub-samples. The ACFs are shown in Fig. 19 and the corresponding ACF exponent $\eta$ is given in Col. 4 of Table 4.

13 The clumpy clustering of the young stars is tightly connected to their self-similar distribution; The more fractal a distribution is, i.e., with lower $D_2$ values, the more prominent sub-clustering behavior its stars have, i.e., higher ACF values (Gouliermis, Hony & Klessen 2014).

14 For a constant star formation rate, a 100 Myr old blue population and a disk of $\sim 10$ Gyr, the young part in the fainter magnitude bins does not exceed $\sim 1\%$ of the old part.
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4.3.2 Pair and Minimum Spanning Tree Separations

According to the results above, since younger stars show a more clumpy clustering behavior than the older, they should be also typically confined in smaller length-scales. We confirm this result with the application of two measurements, the probability distribution function (PDF) of pair separations of all stars in each sub-sample, and the corresponding minimum spanning tree of these stars.

We calculate the PDF of pair separations for all stars in each sample, as the sum of the probability functions of the stars, determined by the number of pair separations that fall in given separation bins around each star (e.g., Cartwright & Whitworth 2004):

\[ p(r_j) = \sum_{i=1}^{N} \frac{2N_{ij}}{N(N-1)d^2} dr. \]  

In this function \( N \) is the total number of stars, \( N_{ij} \) is the number of pair separations that fall in the separation bin centered on \( r_j \), \( r_i \) and \( dr \) is the width of each separation bin. The probability that the projected separation between two randomly chosen stars is in the interval \( (r, r + dr) \) is given by \( p(r)dr \). The constructed pair separations PDFs of young stars show different maxima for different magnitude ranges. Specifically, the PDF becomes systematically more truncated toward smaller separations and its peak occurs at smaller length-scales for stars in younger sub-samples. The median of each derived PDF is given in seconds of arc and kpc in Cols. 5 and 6 of Table 4 respectively. While the RMS dispersion in each PDF is large enough to cover the median of its neighboring PDFs, these values show a clear trend of the typical pair separation of stars toward smaller lengths for higher stellar brightness and age.

The minimum spanning tree (MST), a construct from graph theory, is defined as the unique set of straight lines, called “edges”, connecting a given set of points without closed loops, such that the sum of the edge lengths is the minimum (e.g., Kruskal 1956; Prim 1957). This method has been broadly used, along with other techniques, in clustering analysis of stellar samples (see, e.g. Schmeja 2011, for a review). Here, we construct the PDF of the edge-lengths among stars in every sub-sample, as we did for their pair separations, as defined in Eq. (4). In this measurement instead of the pair separations among all stars, we consider the separations (edge-lengths) of each star from its nearest neighbors, as determined by the MST. The median of the MST edge-lengths for each stellar sample is given in seconds of arc and parsecs in Cols. 7 and 8, respectively, of Table 4. Again, there is a clear trend of younger stars being systematically clumped in smaller scales as defined by the MST.

The similarity of the scaling relations over stellar age found with both methods is demonstrated in Fig. 20, where the corresponding medians are plotted in respect to the limiting magnitude and the equivalent age limit of each stellar sub-sample. Typical uncertainties for the medians, corresponding to a few %, are also shown in Fig. 20. They are calculated from the values correspond-

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**Table 4.** Measures of the clustering behavior of young stars in eight distinct magnitude ranges, corresponding roughly to different stellar age intervals between \( \sim 4 \) and 110 Myr. Magnitudes, ages, and masses listed in the first three columns correspond to the upper limits of each interval. The results of three different measurements are given: The ACF slope \( \eta \) (Col. 4); The median of all pair separations among stars in every sub-sample (Cols. 5 and 6); The median MST edge-length among stars in each sub-sample (Cols. 7 and 8). All measures indicate that the typical length-scale of stellar clustering depends on the brightness (age) range of the stars, with the brighter (younger) being concentrated in systematically smaller scales.

| \( m_{136} \) limit (Myr) | Age (Myr) \( (M_\odot) \) | \( \eta \) | pair sep. median (arcsec) | MST median (arcsec) | MST median (kpc) |
|-------------------------|--------------------------|--------|--------------------------|---------------------|------------------|
| 22.97                   | 32                       | 8.8    | \(-0.693\pm0.004\) | 77.3                | 1.99             | 1.41             | 36.4             |
| 23.62                   | 40                       | 7.7    | \(-0.470\pm0.007\) | 80.2                | 2.06             | 1.85             | 47.6             |
| 24.06                   | 50                       | 6.9    | \(-0.321\pm0.009\) | 80.0                | 2.06             | 1.89             | 48.7             |
| 24.41                   | 63                       | 6.2    | \(-0.289\pm0.011\) | 82.4                | 2.12             | 2.04             | 52.4             |
| 24.72                   | 71                       | 5.9    | \(-0.232\pm0.012\) | 81.9                | 2.11             | 2.19             | 56.4             |
| 25.01                   | 89                       | 5.3    | \(-0.233\pm0.015\) | 84.5                | 2.17             | 2.33             | 60.0             |
| 25.30                   | 100                      | 5.1    | \(-0.225\pm0.016\) | 87.9                | 2.26             | 2.33             | 59.8             |
| 25.91                   | 112                      | 4.9    | \(-0.204\pm0.031\) | 104.9               | 2.70             | 2.96             | 76.0             |

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**Figure 19.** The two-point correlation (autocorrelation) function constructed for eight selected equally-numbered sub-groups of our sample of blue stars, corresponding to different magnitude intervals. In the plot legend the limit of each magnitude bin in the F336W filter is given next to the corresponding color line.
Figure 20. Relation between the medians of (a) the pair stellar separations and (b) the MST edge-lengths and the limiting magnitude (or age) for stars in eight selected brightness ranges (corresponding to different age intervals). These plots demonstrate that there is a dependence of the clustering length-scales of stars to their age. While both measurements show the same trend of scale with age, the difference in the results between the pair separations and MST edge-lengths, which reflects the difference of almost two orders of magnitude between the corresponding medians, lies on the fact that while the PDF of the pair separations is constructed from the whole extent of separations up to the highest covered length-scales, that of the MST edge-length considers only the shortest edges between each star and its neighbors. Under these circumstances, the medians given in Table 4 for the MST edge-lengths represent the most likely typical scale of clustering for each stellar sample, while those for the pair separations correspond to the complete extent where stars in each magnitude bin are located over the whole observed field of view of the galaxy. The pair separations medians provide, thus, the “inter-structure” scaling length of stars at different ages, beyond the “typical” stellar clusterings, determined by the MST.

5 DISCUSSION

The multi-scale sample of young stellar structures detected in NGC 6503 suggests that star formation over a period of \( \sim 100 \) Myr takes place within the typical scale of few hundred parsecs (Sect. 3.1) in structures, which are themselves hierarchical as shown from their dendrogram (Sects. 3.2). The self-similar distribution of all young stars across the galaxy, demonstrated by their ACF, implies that hierarchical structure in the stellar distribution occurs not only inside but also beyond the borders of these structures. Therefore, one can conclude that all star-forming structures are connected to each other in a hierarchical fashion through their coexistence in the star-forming ring of the galaxy (identified in Sect. 4.1).

Hierarchical structural morphology in stellar ensembles has been found by Elmegreen et al. (2014) in the UV images of 12 LEGUS galaxies (NGC 6503 not included in the sample) over a length-scale range of \( \sim 1 - 200 \) pc. The power-law length-scale correlations in the size and flux distribution functions of nested star-forming regions suggest that hierarchically structured star-forming regions with sizes of few hundred parsecs represent common unit structures. This is in line with our findings concerning the typical length-scale of the identified young stellar structures in NGC 6503. Hierarchical structure was identified by Elmegreen et al. (2014) only inside and not among the different star-forming patches, in both large spiral galaxies and low surface brightness dwarfs, as well as in starburst dwarfs or H II galaxies, which are dominated by one or two large star-forming regions. In the case of NGC 6503 self-similar stellar distributions are also detected among the star-forming structures, suggesting that they are all connected to each other in a hierarchical fashion within the dominant star-forming ring of the galaxy.

The observed hierarchical structure of star-forming complexes in NGC 6503 and other galaxies is consistent with the model where star formation is regulated by turbulence, via, e.g., turbulent fragmentation, i.e., gas compressions that form successively smaller clouds inside larger ones (Vázquez-Semadeni et al. 2009). Such processes form a similar hierarchy of young stars, with a likely secondary correlation for star age, making larger regions older in proportion to the turbulent crossing time (Efremov & Elmegreen 1998; de la Fuente Marcos & de la Fuente Marcos 2009). Our findings for stars in different evolutionary stages confirm this time-length correlation, with younger stars being confined in smaller length-scale than the older. The hierarchy in stellar structures has an upper limit in size beyond which separate regions form independently (Elmegreen et al. 2014). In the case of NGC 6503 this limit is set by the size of the star-forming ring of the galaxy. This is consistent with the observation that the two-point correlation for stars decreases as a power law with increasing scale up to about 2.5 kpc.

The peak observed in the young stellar surface density profile of Fig. 16 corresponds to the star-forming ring of the galaxy, and thus its wings define the radial limits of the ring. These limits sug-
gest an inner radius of $\sim 1$ kpc, and an outer radius of $\sim 2.5$ kpc. We derive from the stellar rotational curve of NGC 6503 (Bottema 1989) the projected asymptotic circular velocity for these radii to be of the order of $\sim 110$ km s$^{-1}$ at the outer ring radius and $\sim 80$ km s$^{-1}$ at the inner, corresponding to rotational velocities $\Omega_{\text{out}} \simeq 1.4 \times 10^{-15}$ rad s$^{-1}$ and $\Omega_{\text{in}} \simeq 2.6 \times 10^{-15}$ rad s$^{-1}$ respectively. The rotational velocity difference between the edges of the ring of about $10^{-15}$ rad s$^{-1}$ is comparable to the pattern speed of Milky Way’s spiral arms (Bissantz et al. 2003). It most probably introduces a shear that should influence strongly the star formation process across it.

This shear within $\sim 100$ Myr, which is the time covered by our observed young stellar sample, would have stretched a single compact object inwards, providing that this object is dense enough to survive. The surface stellar density maps of Figs. 7 and 8 suggest an alignment of the detected young stellar structures across the ring of NGC 6503. This trend is visualized by the MST of all structures found at significance levels $\geq 3\sigma$, shown in Fig. 21. In this figure the connecting edge-lengths show an alignment that both follows and traverses the ring closer to its inner part. Considering that the ring has performed several rotations during the last $100$ Myr\textsuperscript{15} it stands to reason that shear functions more as a supportive rather than destructive factor to star formation. The observed clustering of young stars in NGC 6503 can be thus interpreted as being induced by turbulence, the driving source for which is probably gravitational instabilities induced by shear, with larger structures becoming spiral-like (flocculent) because of their longer dynamical time-scales in comparison to the shear time (e.g., Elmegreen 2011).

6 SUMMARY

We present a detailed clustering analysis of the young blue stellar population identified with LEGUS across the star-forming ring galaxy NGC 6503. We construct stellar surface density maps and apply a contour-based analysis technique to identify the stellar complexes population of the galaxy. We identify 244 distinct structures at various stellar density (significance) levels. We organize these structures into 82 separate groups, according to their association with a single low-density parental structure. These groups are arranged into eight families of structures, corresponding to eight super-structures, determined by the $1\sigma$ isopleths. Three of these super-complexes contain 95% of the structures. The hierarchical classification of structures into groups and families, according to their membership to larger and sparser stellar constellations, is illustrated by their dendrogram, or structure tree.

We determine structural parameters, i.e., size, stellar density and total brightness, for each structure. The sizes of the detected systems average around $\sim 130(\pm 40)$ pc, a length-scale comparable to but still larger than $\sim 80$ pc, the scale which is discussed as a characteristic galactic scale for star formation (see, e.g., Gouliermis 2011, and references therein). About 60% of the observed young stellar population is found to belong to one of the $1\sigma$ super-structures, suggesting that $\sim 40\%$ of the young blue stars is distributed in an “unclustered” fashion inside the star-forming ring. The first statistically significant sample of stellar concentrations, detected at the $3\sigma$ density level, accounts for only 34% of the total observed young stellar content, increasing the remaining fraction of “un-clustered” young stars to 66%.

On average, the size (and its dispersion), the total UV brightness, and the fraction of included young stars show a dependence on the density level where the corresponding structures are detected. The fraction of UV light included in the structures over the total observed stellar UV emission shows a steeper dependence on density than that of the stellar fraction. This difference, which becomes more important for higher densities, suggests that compact stellar structures, identified at higher density levels, encompass on average the UV-brightest stars in the galaxy.

We identify a power-law mass-size relation, determined by the correlation between radius and number of stars. The exponent of this relation for all structures is less than $2$ as expected for a fractal distribution of stars (e.g., Elmegreen & Falgarone 1996). Our findings suggest a dependence of the power-law index of the mass-size relation to the detection threshold, in agreement with previous studies in the Large Magellanic Cloud (Gouliermis et al. 2003) and M 33 (Bastian et al. 2007). A correlation between the volume stellar density and the size of the structures was also found. This relation for the systems detected at 5 and $6\sigma$ levels resembles Larson’s third relation (Larson 1981) and becomes steeper for higher-density structures, possibly reflecting their steeper density profiles. Since the volume density scales inversely with size, the identified correlation implies a dependence of the parameters determination for the structures on the detection criteria, in a similar manner as for GMCs.

The radial surface density profile of the young stars in NGC 6503 shows a very strong central peak and drops rapidly within a radial distance of $\sim 350$ pc, in agreement with the line-of-site velocity dispersion profile, which shows a sharp decrease within $300$ pc (Bottema 1989; Bottema & Gerritsen 1997). The stellar density profile rises again and reaches its peak at a distance of $\sim 1.9$ kpc, after which drops again exponentially. This profile, demonstrating a circumnuclear deficiency in young stars, agrees with both an inner-truncated disk and a ring being formed by a bar (Puglioni et al. 2010). The confinement, though, of all young stellar structures in a ring-shaped alignment favors the ring-formation scenario, as earlier suggested (Freeland et al. 2010).

We characterize the hierarchy in the global clustering behavior of young stars across NGC 6503 with the autocorrelation function (ACF). We find that the ACFs of these stars shows the typical features of a self-similar stellar distribution with a two-dimensional fractal dimension $D_2 = 1.7$. The observed hierarchy in young stellar clustering extends monotonically across the complete measurable dynamic range in length-scales of two orders of magnitude ($\sim 20$ - $2500$ pc). The self-similar distribution of young stars across NGC 6503 is consistent with the hierarchical morphology of the star-forming complexes of the galaxy, and the mass-size relation of the structures, which fits to the expectations for fractal clusterings.

We investigate the clustering behavior of young stars at different evolutionary stages with the ACF of stars in different magnitude intervals. We find that younger (brighter) stars are more clustered than the older. The ACF exponents of the younger stars are systematically higher, corresponding to smaller fractal dimensions and more clumpy distributions, than those for the older stars. This analysis shows that the time-scale, where a significant change from a more clustered to a more distributed assembling of stars takes place, is about 60 Myr. A similar trend was found for the Small and Large Magellanic Clouds, with a time-scale for sub-structure evolution towards a uniform distribution of $\sim 75$ and 175 Myr, respectively. We determine structural parameters, i.e., size, stellar density and total brightness, for each structure. The sizes of the detected systems average around $\sim 130(\pm 40)$ pc, a length-scale comparable to but still larger than $\sim 80$ pc, the scale which is discussed as a characteristic galactic scale for star formation (see, e.g., Gouliermis 2011, and references therein). About 60% of the observed young stellar population is found to belong to one of the $1\sigma$ super-structures, suggesting that $\sim 40\%$ of the young blue stars is distributed in an “unclustered” fashion inside the star-forming ring. The first statistically significant sample of stellar concentrations, detected at the $3\sigma$ density level, accounts for only 34% of the total observed young stellar content, increasing the remaining fraction of “un-clustered” young stars to 66%.

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spectively (Gieles et al. 2008; Bastian et al. 2009). Larger scale occu-
pancy by stars of decreasing luminosity is previously discussed
also for the spiral NGC 1313 and dwarf irregular IC 2574 (Pellerin
et al. 2007, 2012). These results are conﬁrmed with the pair sepa-
rations and Minimum Spanning Tree edge-lengths probability dis-
tributions of the stars in the various magnitude (age) ranges.

Based on kinematic arguments we assess that hierarchy in the
young stellar clustering in NGC 6503 is probably induced by tur-
bulence, driven by shear in the ring of the galaxy. With this mech-
nism large stellar structures become ﬂocculent because their dy-
namical time-scales exceed the shear time.

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16 TOPCAT is available at the permalink http://www.starlink.ac.uk/topcat/
17 Accessible at http://adswww.harvard.edu/ and http://cdsads.u-strasbg.fr/
