Preconditioning of Interplanetary Space Due to Transient CME Disturbances

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Abstract

Interplanetary space is characteristically structured mainly by high-speed solar wind streams emanating from coronal holes and transient disturbances such as coronal mass ejections (CMEs). While high-speed solar wind streams pose a continuous outflow, CMEs abruptly disrupt the rather steady structure, causing large deviations from the quiet solar wind conditions. For the first time, we give a quantification of the duration of disturbed conditions (preconditioning) for interplanetary space caused by CMEs. To this aim, we investigate the plasma speed component of the solar wind and the impact of in situ detected interplanetary CMEs (ICMEs), compared to different background solar wind models (ESWF, WSA, persistence model) for the time range 2011–2015. We quantify in terms of standard error measures the deviations between modeled background solar wind speed and observed solar wind speed. Using the mean absolute error, we obtain an average deviation for quiet solar activity within a range of 75.1–83.1 km s⁻¹. Compared to this baseline level, periods within the ICME interval showed an increase of 18%–32% above the expected background, and the period of two days after the ICME displayed an increase of 9%–24%. We obtain a total duration of enhanced deviations over about three and up to six days after the ICME start, which is much longer than the average duration of an ICME disturbance itself (∼1.3 days), concluding that interplanetary space needs ∼2–5 days to recover from the impact of ICMEs. The obtained results have strong implications for studying CME propagation behavior and also for space weather forecasting.

Key words: solar–terrestrial relations – solar wind – Sun: coronal mass ejections (CMEs) – Sun: heliosphere

1. Introduction

The solar wind during quiet-Sun conditions consists of a continuous outflow of slow solar wind streams from the low-latitude streamer belt and fast solar wind stemming from coronal holes. This pattern of alternate slow and fast solar wind streams repeats itself with each solar rotation. It is disrupted by transient events, such as coronal mass ejections (CMEs), which present dynamic regions of reduced plasma density, enhanced speed, and radially stretched magnetic field, therefore causing strong changes in the solar wind properties on much shorter timescales (see, e.g., Schwenn 2006). Such disturbed conditions of interplanetary space may be labeled as “preconditioning” for subsequent CMEs expelled into these regions of interplanetary space. As CMEs evolve from the Sun, they are initially governed by the Lorentz force and in their later evolution phase by the aerodynamic drag force (see, e.g., Cargill et al. 1996). The drag force becomes dominant as the CME propagates into interplanetary space and is usually reflected in the deceleration of CMEs that are faster than the ambient solar wind speed and acceleration of CMEs that are slower than the solar wind (Gopalswamy et al. 2000; Vršnak 2001).

Keeping in mind that the average occurrence frequency of CMEs is about 0.5 per day during the solar minimum phase and 2–3 per day during solar maximum (Cyr et al. 2000), where successive CMEs are often launched from the same active region, we may assume that a continuous, i.e., nondisturbed, solar wind flow toward a certain direction prevails only during times of low solar activity. This is also reflected in the performance of background solar wind models. Reiss et al. (2016) showed for different (semi-)empirical background solar wind models that the correlation coefficient between modeled and observed solar wind speed at Earth is higher for times without CMEs (see also Kohutova et al. 2016). Even sophisticated numerical MHD background solar wind models tend to fail during times of enhanced solar activity, while they are well matching the actual observation during times of low activity (e.g., Jian et al. 2011; Gressl et al. 2014).

The conditions of interplanetary space, i.e., the ambient magnetic field and plasma flow, in which CMEs are embedded, affect the CME transit times and their geo-effectiveness. Preconditioning of interplanetary space due to transient disturbances has immediate consequences for the propagation behavior of CMEs, as was, e.g., impressively revealed by the event from 2012 July 23. This CME was observed to be superfast and traversed the distance from Sun to STEREO-A in less than 21 hr, as a result of the changed conditions of interplanetary space most probably caused by the CME launched 3.5 days earlier from the same active region (Liu et al. 2014; Temmer & Nitta 2015). If Earth-directed, this event would have been one of the most extreme space weather events, with an estimated Dst > 1100 nT (Baker et al. 2013). Thus, the question arises, how long does the interplanetary space need to recover from CME disruptions?

In this study we will, for the first time, quantify the duration of preconditioned phases of interplanetary space as a consequence of preceding CMEs. To this aim, we apply background solar wind models and compare the modeled solar wind speed to actual observations for times with and without CMEs. We investigate the time range 2011–2015, which is favorable since solar cycle 24 is rather weak and we may expect extended periods of undisturbed solar wind flow. In situ measurements show for time ranges influenced by interplanetary CME (ICME) disturbances on average plasma speeds clearly above the background level, in addition to significant signatures of compressed sheath plasma behind shocks and a smooth magnetic field inside the ejecta clouds.
(e.g., Schwenn 2006). We will show that these preconditioned periods prevail for much longer time ranges than the duration of the ICME disturbance itself. In fact, we find that the impact of CMEs affects interplanetary space over periods from 3 to 6 days. This has important implications for CME propagation models and space weather forecasting algorithms, as precursor events over a period of up to 6 days may need to be taken into account.

2. Data and Methods

We investigate the effect of ICMEs on the interplanetary solar wind speed profile for the time span 2011–2015. We use level-2 solar wind data (solar wind bulk speed [km s⁻¹]) from the ACE/SWEPAM instrument located at L1 (McComas et al. 1998). The 64 s resolution solar wind data were linearly interpolated to a 4-minute time resolution. The arrival times and mean impact speeds of the ICMEs (shock-sheath, magnetic cloud structure) are extracted from the Richardson and Cane list (“R&C List”; Richardson & Cane 2010).

To simulate the continuous, quiet-Sun background solar wind speed (without CME disturbances), we use the empirical solar wind forecast (ESWF) and the Wang–Sheeley–Arge (WSA) model. The ESWF and WSA model results are given for Earth distance and cover a cadence of 6 and 4 hr, respectively. For comparison, all data used for the present study are linearly interpolated onto a 6 hr time grid. The ESWF models the quiet solar wind speed at 1 au based on the fractional area of coronal holes detected in EUV filtergrams at the central meridian of the Sun covering a slice of 15° longitudinal width corresponding to the solar rotation within approximately 1 day (see Vršnak et al. 2007; Rotter et al. 2012). The coronal hole areas are extracted using an intensity-based threshold technique applied on SDO/AIA image data in the 19.3 nm band with a time cadence of 6 hr (see Reiss et al. 2016). In contrast, the WSA model empirically relates the solar wind speed to the flux-tube expansion factor of open magnetic field lines (Arge & Pizzo 2000; Arge et al. 2003), which are computed from the potential-field source-surface model (Altschuler & Newkirk 1969) applied on magnetograms from the Global Oscillation Network Group. For more details on the application of this model to derive the background solar wind speed at 1 au, we refer to Nikolic et al. (2014). As a cross-check, we use a 27-day persistence model (PS27), which uses the solar wind speed measured on day d to model the speed on day d + 27. Reiss et al. (2016) performed a systematic comparison of the three models with the in situ measured solar wind speed at 1 au for the time range 2011–2014 (excluding ICME periods) and found that all three model performances give reasonable results enabling us to reproduce the quiet solar wind conditions.

As described above, the models used for this study are capable of reproducing large-scale features of the continuously emanated solar wind, but do not take into account transient events such as CMEs. By comparing the in situ measured solar wind speed data at L1 with the model results, we can therefore study the deviations caused by Earth-impacting CMEs. As a measure for the differences between actual in situ observations and background solar wind, we compute the (a) mean error (ME; arithmetic mean between observed and model results); (b) mean absolute error (MAE; arithmetic mean of absolute differences between observations and model results); and (c) root mean square error (RMSE; mean squared difference between observations and model results). The ME indicates disproportionate shifts toward positive or negative differences, and the RMSE is computed by squaring the differences, hence giving more weight to large errors. The MAE can be seen as the most realistic estimate of the uncertainties between the modeled background solar wind speed and the actually measured solar wind speed.

We would like to note that all algorithms used to conduct the present analysis, together with the scripts of the figures, are openly accessible online at https://bitbucket.org/reissmar/icme-preconditioning.

3. Results

Figure 1 shows for the entire time range under study (2011–2015) the derived differences between measured solar wind speed (ACE/SWEPAM) and modeled solar wind speed. In total, there are 136 ICMEs listed during 2011–2015, whose duration (ICME shock arrival until the end of magnetic structure) is marked with green vertical bars. A total of 22 of the 136 ICMEs arrive simultaneously with high-speed solar wind streams. The mean impact speed of the ICMEs is derived with 430 ± 76 km s⁻¹ (minimum and maximum speed is 290 and 680 km s⁻¹, respectively). Inspecting the profiles, we find that CMEs often lead to prolonged speed differences after their transit. In general, PS27 reveals smaller deviations for times without ICMEs compared to ESWF and WSA.

Figure 2 gives the errors (ME, MAE, RMSE) describing the deviation between modeled background solar wind speed and observed solar wind speed profiles calculated for three time intervals: (1) time interval tICME covers the entire ICME time range, starting with the arrival of the shock-sheath structure until the end of the magnetic cloud; (2) tICME covers all data points within two days starting after the end of the magnetic cloud; (3) tICME covers all data points within two days starting after the end of the magnetic cloud; (4) tICME covers all data points within two days starting after the end of the magnetic cloud.

As presumed, this clearly reveals that all background solar wind models are associated with larger errors during the occurrence of transient events. The MAE calculated for the different models over the interval tICME lies in the range of 75.1–83.1 km s⁻¹, which we refer to as the “baseline level,” indicating the undisturbed conditions. Compared to this baseline level, the MAE is enhanced during tICME by 18%–32% and during tICME by 9%–24%. A similar outcome is derived for the RMSE (see Table 1 for all derived results using different models and error measures). From this we conclude that even within two days after the magnetic structure of the ICME has passed the spacecraft, the solar wind speed did not go back to the baseline level. This is analyzed in more detail in the following.

We first illustrate the differences between the background solar wind speed and the ICME mean impact speed (see column (i) in the “R&C List”), for which we extract for each ICME the period covering 2.5 days before and 10 days after the arrival of the ICME (t = 0). Figure 3 shows for each of the 136 ICME events the speed differences between observations and model as a function of time (with t = 0 referring to the arrival time of the ICME shock-sheath structure), stacked in order by the ICME mean impact speed (y-axis top to bottom).
fastest to slowest events). With the start of the ICME impact at the spacecraft, the deviation between modeled and measured speed increases, lasting for at least three days. From this, no clear relation is obtained between ICME mean impact speed and duration of disturbed conditions.

While Figure 3 provides a visual assessment of the deviations between modeled and measured speed, Figure 4 provides the quantification of the required time for the interplanetary medium to recover from transient events. To this end, Figure 4 shows the computed error measures ME,
MAE, and RMSE (deviations between observed and modeled solar wind speed) as calculated from the different models (ESWF, WSA, PS27) covering the period from 2.5 days before to 10 days after the arrival of the ICME ($t = 0$). We give the results for four different ICME categories, which correspond to different ICME mean impact speeds. To ensure a sufficiently large sample size, the four categories cover 136 ICMEs with $v > 300$ km s$^{-1}$, 122 ICMEs with $v > 350$ km s$^{-1}$, 78 ICMEs with $v > 400$ km s$^{-1}$, and 46 ICMEs with $v > 450$ km s$^{-1}$.

The reference time zero is the shock arrival time of each ICME at the in situ spacecraft. Hence, the derived similarities in the error profiles of the four speed categories reflect a characteristic behavior of the evolution of the disturbances. As an indication of the undisturbed conditions, the calculated deviations outside recorded ICME disturbances ($t_{\text{out}}$), hence the baseline level for the MAE results, are given by blue dotted horizontal lines.

### Table 1

| Model | Error | $t_{\text{ICME}}$ | $t_{1.2d}$ | $t_{\text{out}}$ |
|-------|-------|-------------------|-----------|-----------------|
| ESWF  | ME    | 55.9              | 44.1      | 16.0            |
| ESWF  | MAE   | 97.8              | 90.4      | 83.1            |
| ESWF  | RMSE  | 125.5             | 117.0     | 108.2           |
| WSA   | ME    | 87.2              | 78.6      | 43.7            |
| WSA   | MAE   | 101.5             | 95.5      | 77.1            |
| WSA   | RMSE  | 130.1             | 123.4     | 103.4           |
| PS27  | ME    | 37.0              | 34.7      | 1.3             |
| PS27  | MAE   | 89.4              | 92.3      | 75.1            |
| PS27  | RMSE  | 120.8             | 121.2     | 101.7           |

**Figure 3.** Speed differences (color-coded) for each of the recorded ICMEs stacked by the ICME impact speed as a function of time. The differences are calculated from different background solar wind models (top to bottom: ESWF, WSA, PS27). The reference time zero refers to the shock arrival of the ICME.
ICME shock-sheath and magnetic structures are marked by vertical red dot-dashed and blue/black dashed lines, respectively. The listed events under study have a mean duration for the entire ICME, $t_{\text{ICME}}$, of $1.27 \pm 0.67$ days and a mean duration for the magnetic structure of $0.96 \pm 0.53$ days (calculated from the “R&C List”; Richardson & Cane 2010). Different colors for the profiles of the calculated error measures indicate the four ICME mean impact speed categories (see legend for speed and sample size). The blue dotted horizontal line is the deviation calculated outside recorded ICME disturbances ($t_{\text{out}}$). Focusing on the results from the MAE (middle panels), we mark the mean ICME duration and the recovery time (red arrows), hence the time until the increased deviations drop back to the baseline level.

ICME shock-sheath and magnetic structures are marked by vertical red dot-dashed and blue/black dashed lines, respectively. The listed events under study have a mean duration for the entire ICME, $t_{\text{ICME}}$, of $1.27 \pm 0.67$ days and a mean duration for the magnetic structure of $0.96 \pm 0.53$ days (calculated from the “R&C List”). Inspecting the error profiles for the entire sample ($v > 300 \text{ km s}^{-1}$), all models reveal the largest deviations during the passage time of the ICME shock-sheath structure, followed by a rather linear decrease of the error profile during the passage of the magnetic cloud structure. From ESWF and WSA models and for each ICME category, we derive that all computed error measures continue to be higher than the baseline level and only start to decrease at about $t_0 + 3$ days, lasting until about $t_0 + 6$ days (for faster events with $v > 400 \text{ km s}^{-1}$). This recovery time is much longer than the average duration of the entire disturbance. Independent from the ICME category, measure errors, and models used, a short-term enhancement in the deviations around 4–5 days is revealed. The higher the mean impact speed of the ICME event, the larger the derived values for the ME, MAE, and RMSE. We note that with the smaller subset for faster ICMEs, larger fluctuations in the errors before the shock arrival are derived, which is related to the smaller statistics.

### 4. Discussion and Conclusions

In this study, we evaluate the impact of ICMEs on the solar wind conditions in interplanetary space for the time range 2011–2015. To simulate the quiet solar wind speed, we use different (semi-empirical models such as the ESWF, WSA, and PS27. For periods affected by ICMEs we find clear deviations that last at least three days (up to six days) after the
ICME start. This is longer than the duration of the ICME disturbance itself (for our sample on average $\sim1.3$ days) as it passes the in situ spacecraft. Prolonged geomagnetic storms and long-duration magnetic cloud structures as a consequence of ICME flank hits are found to last $\sim$1.5–2 days (Marubashi & Lepping 2007; Möstl et al. 2010). Although the investigated approach introduces uncertainties due to the analysis of stacked events, our results for the first time quantify over which time interval ICME disturbances affect/precondition interplanetary space (total disturbance period is $\sim$3–6 days) and how long interplanetary space may need to recover from that (total disturbance period minus ICME duration is $\sim$2–5 days). With these results, we can confirm conjectured scenarios explaining why the CME from 2012 July 23 became superfast (Liu et al. 2014; Temmer & Nitta 2015). Besides the solar wind speed, we presume that also other solar wind properties, such as the magnetic field and the density, are affected over a similar time span.

There is some indication that faster ICMEs might disrupt the interplanetary medium for a longer time span compared to slow ICMEs. However, the derived fluctuations are found to be larger for higher ICME impact speeds. We stress that the time range under study covers the weak solar cycle 24; hence, very fast ICME events are actually missing in this sample. We also note that a fraction of about 15% of the listed ICMEs are related to enhanced solar wind speed streams emanating from coronal holes (for the relation of stealth ICMEs to coronal holes see, e.g., D’Huys et al. 2014). For those events we derived no significant deviations between modeled and observed solar wind speed.

The question remains, what causes such specific long-lasting preconditioning periods? (1) The time interval is comparable to the average solar wind transit time from Sun to Earth and might reflect the time span that the entire system, i.e., the interplanetary medium, needs to recover from the propagating disturbance. (2) The occurrence of multiple ICMEs, overlapping with each other, might cause prolonged disturbances. We cross-checked this possibility for a subset of 57 events, for which all ICMEs that had other ICMEs surrounding them with $\sim$2.5 days were removed, and performed the analysis again. We still found the same trend of prolonged preconditioning. This may indicate that (3) other kinds of disturbances, or nonlisted ICMEs, traveling in the wake of the listed ICMEs might cause this effect (the derived short-term increase in the deviations around 4–5 days might support this scenario). Aftermath regions of ICMEs without clear rotating magnetic field structures could be due to reconnection of the magnetic structure with the interplanetary magnetic field (Dasso et al. 2007). Peaks in plasma density and speed observed close to the trailing edge of magnetic flux ropes were interpreted as compression caused by high-speed solar wind streams running into the expanding ICME (Rodríguez et al. 2016).

The obtained results have strong implications for studying CME propagation behavior and also for space weather forecasting, as they provide quantitative information on the reliability of operational forecast models during ICME disturbances.

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