Measurement of Coherent Undulator Radiation of Compact Terahertz Radiation Source at Kyoto University

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Abstract
An accelerator based compact terahertz radiation source has been successfully developed at Institute of Advanced Energy, Kyoto University. The system is very compact with the total length of 3.7 m. The radiations are generated by injecting a short electron bunch to a 0.7 m long 10-period planar Halbach type undulator. The radiations have a high degree of phase coherence or called coherent undulator radiation. The frequencies are tunable in the range from 0.17 to 0.55 THz. The electron source is a 1.6 cells BNL type photocathode RF gun which provides the beam with energy of 4.6 MeV. The beams are compressed in longitudinal by a magnetic chicane bunch compressor. The first light was generated in Aug 2016. The results of the measurement of the light properties are presented in this paper.

Keywords
Undulator radiation, Coherent radiation, Terahertz

Introduction
Over past few decades, the studies based on Terahertz (THz) frequency electromagnetic radiation or terahertz science gain great interests in various fields from fundamental researches to industrial applications. The frequency of THz radiation is in the far infrared spectral range (approximately 0.1 to 10 THz) which corresponds to the intermolecular vibration [1]. THz radiation allows us to extend the study of the molecular dynamics beyond conventional infrared radiation techniques leading to new information or methods of a material analysis, a medical diagnostic, a non-destructive inspection or a security, etc. However, the THz radiation is difficult to generate by conventional devices. Therefore, the development of the THz source is one of the main challenges in THz science, especially in the high power range.

An accelerator-based THz radiation source is a kind of sources that can provide significantly high power, tunable frequency and narrow bandwidth. There are many accelerator-based THz radiation sources operating worldwide, such as, at JLAB [2], DESY [3], SLAC [4]. However, these systems are large-scale facilities that require high cost for the construction and the operation. Over past decade, a small-scale accelerator system has been proven that it has the potential to achieve high power and continuous tunability in the THz radiation regime. The system employs bunched electron beams whose bunch length comparable to the radiation wavelength injecting to a short undulator. The radiation from the undulator has a high degree of phase coherence or called Coherent Undulator Radiation (CUR). There are a number of systems that have been developed, for examples the compact systems at ENEA (0.4-0.7 THz) [5] and Peking Univ. (0.24-0.42 THz) [6] etc.

We also have been developed accelerator-based compact THz radiation source at Institute of Advanced Energy, Kyoto University. The system is able to produce a high power tunable THz CUR by injecting a short electron bunch to a short planar undulator. The system is compact with the total length of 3.7 m and shares the building and the RF power source with Kyoto University Free-Electron Laser (KU-FEL) [7]. The first photoelectron beam was generated from a 1.6 cells S-band BNL-type photocathode RF gun in April 2015 [8]. The mag-
netic chicane bunch compressor was installed in March 2016 and the compressed electron bunch characteristic was investigated by the coherent transition radiation technique in April 2016 [9]. The undulator was installed in July 2016. The first THz CUR was successfully generated in August 2016. The characteristic of the radiation were measured by two small-size uncooled detectors and the spectra were analyzed by an in-air Michelson interferometer. This paper presents the details, results and discussions of the measurement of the first THz CUR. However, the absolute values of the intensity were not calculated because the precise calibration of the measurement system is still underway.

Coherent Undulator Radiation

Undulator is a device which generates a periodic magnetic field to drive the electron in a sinusoidal trajectory. During electrons travel through an undulator, electrons emit the radiations in the forward direction as shown in Figure 1 and the radiation wavelength (λ) can be determined by

$$\lambda = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + \gamma^2 \theta^2\right)$$  \hspace{1cm} (1)

Where λ_u is the undulator period length, y is the Lorentz factor, θ is the observation angle from the beam direction, K is the undulator parameter, K = eB_0/mecku m_e is the charge of an electron, B_0 is the peak magnetic field, me is the rest mass of electron and c is the speed of light and k_u is the undulator wave number, 2π/λ_u. The radiations concentrate in a narrow cone with the opening angle of K/γ. The Coherent Undulator Radiation (CUR) can be obtained when the electron bunch length is comparable to the radiation wavelength. The total field of CUR (E_{tot}(ω)) at frequency ω can be determined by a sum of the emitted fields from each electron [10] as

$$E_{tot}(\omega) = N^2 f(\omega) E_z(\omega)$$  \hspace{1cm} (2)

Where N is the number of electrons, f(ω) is the Fourier transform of the electron distribution or called the form factor and E_z is the field of a single electron. The f(ω) in the longitudinal direction can be calculated by

$$f(\omega) = \int s(z) e^{i\omega z} dz$$  \hspace{1cm} (3)

Where s(z) is the normalized electron bunch distribution in longitudinal direction, z is the electron position in longitudinal direction and k is the wave number of the radiation, k = ω/c = 2π/λ. The total intensity of the CUR (I(ω)) equal to the total field square, I(ω) = |E_{tot}(ω)|^2. To obtain a high CUR intensity, we employed the bunch compressor to compress electron bunch longitudinally to a shorter bunch length resulting in the enhancement of the form factor and the CUR intensity. The schematic diagram of the bunch compression in longitudinal is shown in the Figure 2.

This bunch compressor consists of four dipole magnets and requires the bunch with a low energy at the head and a high energy at the tail as shown in Figure 2. When such a bunch passes through the chicane, the higher energy electrons travel in the shorter path and can catch up the low energy electron at the front resulting in a shorter bunch length at the chicane exit. The suitable initial longitudinal phase space correlation of the bunch for the compression can be obtained by adjusting the laser injection phase to the RF gun. The amount of compression is determined by the first order momentum compaction factor (R_56) which is defined as the ds/dy here ds is the total path different and dy is the relative energy.

The spectral intensity of the CUR is not monochromatic due to the finite length of the CUR wave train. The wave train has the number of oscillation same as the number of periods of the undulator (N_u) and the time duration of T = N_u/λ/c. The spectral bandwidth (Δω/ω_0) of the CUR depends on the number of period of the undulator that can be calculated by

$$\frac{\Delta \omega}{\omega_0} = \frac{1}{N_u}$$  \hspace{1cm} (4)

Where ω_0 is the central frequency of the radiation.
Undulator of Compact THz Radiation Source

The undulator of the compact THz radiation source at Kyoto University is a vertical deflection planar Halbach type permanent magnet with the period length of 70 mm and the number of periods of 10. The undulator gap can be manually adjusted with the minimum gap of 30 mm limited by the vacuum chamber. At the gap of 30 mm, the measured peak magnetic field at the center between poles is 0.43 T and the maximum undulator parameter ($K$) is 2.8 corresponding to the lowest radiation frequency of 0.17 THz. The expected spectral bandwidth is 10% for all CUR frequencies. The measurement of the transverse magnetic field at the center between poles at the undulator gap of 30 mm is shown in Figure 3. The calculated undulator parameter as a function of the undulator gap is shown in Figure 4.

Measurement Setup

The beam line of the compact THz radiation source is shown in Figure 5. The electrons were generated by injecting the 266 nm wavelength UV laser [11] to the copper cathode of the 1.6 cells photocathode RF gun. The laser had the repetition rate of 89.25 MHz which was the 1/32 time of the driven RF frequency (2,856 MHz, S-band) for the RF gun and the pulse duration at the FWHM approximately of 6 ps. The RF pulse had the macropulse repetition rate of 2 Hz and the time duration of 2 μs. The number of laser pulse was adjusted from 1 to 4 pulses per macropulse to obtain proper detected signal by detectors. This experiment operated the RF gun at the maximum power of 9 MW that provide the beam with energy around 4.6 MeV and energy spread of 1.3% [9]. The magnetic chicane bunch compressor operated in two conditions which were turn-on and turn-off compression. The DC current of 5.8 A supplied to the chicane dipole magnets for the turn-on compression that provided the $R_{56}$ of -47.1 mm. The solenoid magnet attached in front of the RF-gun, the steering magnets, the small correction coil of the chicane dipole magnets and triplet quadrupole magnets were adjusted to obtain the highest detected signal from the detectors, i.e. highest radiation intensity.

THz radiations from the undulator reflected out of the vacuum chamber by a retractable THz mirror made from a titanium film with the thickness of 50 micron installed in a 6-crosses chamber downstream the undulator. The film was attached on a circular aperture aluminum holder with the aperture diameter of 50 mm arranged in the 45 degree respect to the electron beam direction and 480 mm apart from the middle of the undulator. The fused silica with the aperture diameter 65 mm was employed as a THz window and the distance from the mirror to the window was about 120 mm. Two types of small-size uncooled detectors were used for the THz radiation detection. First, a Schottky Zero-Bias Barrier Diode (SBD) detector (Millitech, DET-05-RPFW0) has the sensitivity of higher than 250 mV/mW [12] in the frequency range from 140 to 220 GHz. The SBD detector was installed inside the rectangular wave guide, therefore, it could pick the signal in one direction of polarization. Second, a pyroelectric detector (PHLUXi, PYD-1-018) equipped with the built-in lens and the visible light filter. The detector has not been calibrated in the THz frequency range. The signals from the detectors were recorded by 1 GHz oscilloscope (Textronics, DPO4104). The bunch charge was measured by a 10 cm cubic out-of-vacuum carbon
Faraday cup placed at the end of the beam line and the signal was read by 200 MHz oscilloscope (Iwatsu, DS-5624A). The lead blocks were put beside the beam line to prevent bremsstrahlung radiation inducing noise in the detector. The beam line control and measurements were performed under the LABVIEW environment. The parameters for the undulator radiation measurement are listed in Table 1.

### Spatial distribution

The spatial distributions of the undulator radiation were measured by scanning the SBD detector in two-dimensions without focusing optics. The scanning plane was normal to the extracted THz beam with the distance approximately of 140 mm from the center of the THz mirror. The detector was mounted on two translational stages (Sigmakoki, SG SP 20-85) and scanned with 4 mm interval in both horizontal and vertical directions. The measurement set up of the spatial distribution of the undulator radiation is shown in Figure 6.

### Total intensity

The total intensity was measured by using a focusing optics to focus whole extracted beam to the pyroelectric detector. A plano-convex polytetrafluoroethylene lens or Teflon lens (Thorlabs, LAT115) with the 3 inch diameter and the focal length of 115 mm was employed and placed next to the THz window. The detector was aligned to the focal point by means of the translation stage. The setup of the total intensity measurement is shown in Figure 7.

### Power spectrum

The setup of an in-air Michelson interferometer for the power spectrum measurement is shown in Figure 8. Two focusing lens which were the 3-inch and 2-inch diameter Teflon lens (Thorlabs, LAT115, LAT75) with the focal length of 115 and 75 mm were utilized for reducing the THz beam size suitable for the interferometer. The fixed and movable mirrors of the interferometer were 2-inch gold-coated flat mirror (Thorlabs, PF20-03-M01). For measuring the interferogram, the movable mirror position was varied from -8 to 8 mm with the step size of 0.1 mm. The total optical path different was 32 mm and the time window of the measurement was 106.7 ps and the frequency resolution, which is the inverse of the time window, was 9.37 GHz.

Two types of the beam splitter which were a 100 micron thick sapphire beam splitter and an inconel coated pellicle beam splitter with the thickness of 2 micron were used. Due to the efficiency drop of the sapphire

| Parameters                                | Values                              |
|-------------------------------------------|-------------------------------------|
| Beam energy                               | 4.6 MeV                             |
| Energy spread                             | 1.3%                                |
| Laser pulse duration                      | 6 ps (FWHM)                         |
| Laser size at cathode (Hor./Ver.)         | 0.4/0.5 mm rms                      |
| Laser distribution (Hor./Ver.)            | Gaussian/Gaussian                   |
| Number of laser pulse per macropulse      | 1 to 4 pulses                       |
| Solenoid magnetic field                   | 163 to 194 mT                       |
| $R_{56}$ of the magnetic chicane           | 0, -47.1 mm                         |
| Undulator gap                             | 30 to 55 mm                         |

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beam splitter at around 0.5 THz, this beam splitter was
used for the measurement at the frequency less than 0.4
THz while the inconel coated beam splitter was used at
the frequency more than 0.4 THz. The combined beams
exiting the interferometer were focused by the 2-inch
Teflon lens with the focal length of 50 mm (Thorlabs,
LAT50). The pyroelectric detector with the built-in lens
was mounted on the translation stages and aligned to the
focusing point of the combined beam.

Measurement Result
Spatial distribution
The spatial distribution measurements were per-
formed at the laser injection phase of 20 deg and the $R_{36}$
of -47.1 mm which provided the best compression con-
dition as studied in [9]. The undulator gap was varied to
30, 35 and 180 which were the minimum gap, the gap at
the highest frequency response of the SBD detector
and the largest gap. These gaps had the undulator pa-
rameter ($K$) of 2.8, 2.3 and 0 and the calculated central
frequencies of the radiation ($\omega_0$) of 0.17, 0.23 and 0.86
THz, respectively. The number and the energy of the la-
sers pulse per macro pulse were 4 pulses and 300 μJ. The
results of the spatial distribution of the undulator radia-
tion are shown in Figure 9. Note that all shown results
were the radiation with the vertical polarization which
is the same direction as the deflection of electron inside
the undulator. The radiation with the horizontal polar-
ization was also measured but the signal was very weak.
As the results, the radiation from the compressed bunch
concentrated in the narrow cone at the beam center and
aligned to the center of the THz mirror (Figure 9a). The
radiation from the uncompressed bunch was weak (Fig-
ure 9d). At the gap of 30 and 35 mm, the calculated open-
ing angle ($K/\gamma$) of the radiation cone corresponding to
the beam energy of 4.6 MeV ($\gamma$=10) were 16 and 13.1
deg that larger than the acceptance angle of the mirror at
around 6 deg. Therefore, only the center part of the radi-
Figure 9: Spatial distributions of the CUR of the compressed bunch at the undulator gap of a) 30 mm; b) 35 mm; c) 180 mm
and the uncompressed bunch at the undulator gap of d) 30 mm.
ation could be reflected by the mirror. But the other part of the radiation might be reflected by the internal side of the 6-cross chamber as clearly seen in Figure 9b. In case of the largest undulator gap at 180 mm (Figure 9c), the CUR could not be observed but the detected radiation was a coherent transition radiation, which was emitted at the surface of THz mirror made of Ti film when electron bunch was injected to the film.

### Total intensity

The measured bunch charge and the radiation intensities dependence on the laser injection phases of the uncompressed and compressed bunch are shown in Figure 10 and Figure 11. The measurement was performed with the number and the energy of the laser pulse per macro-pulse of 4 pulses and 300 μJ, respectively. The measured bunch charge of both compression conditions had almost the same for a given laser injection phase and undulator gap. But the CUR intensity of the compressed bunch was significantly higher than the uncompressed bunch. It means that the chicane successfully compressed the bunch and enhanced the form factor resulting in a high CUR intensity. The peak intensity occurred at the phase around 15 to 25 degrees, where the bunch had suitable initial longitudinal phase space correlated with the compression at the $R_{56}$ of -47.1 mm. The factor of the intensity enhancement reached up to 8.3 at the undulator gap of 50 mm. The CUR from the larger undulator gap had smaller intensity compared to the narrower one, because the form factor of the same bunch became smaller at a higher frequency.

### Power spectrum

The autocorrelation intensity dependencies as a function of path different or interferograms measured by the Michelson interferometer for the undulator gap from 30 to 55 mm are shown in Figure 12. Note that the interferograms are shifted in vertical for easier seeing the lines. The measurement was performed only for the compressed bunch because the signal had strong enough to detect. The undulator radiation spectra calculated by Fast Fourier Transform corresponding to the interferograms are shown in Figure 13. As the result, the central
low energy electron beam of 4.6 MeV. The radiation was coherent and the intensity could be enhanced by the magnetic chicane bunch compressor. For the compressed bunch at the $R_{cm}$ of -47.1 mm, the radiation intensity had the peak at the laser injection phase of 15 to 25 degrees. The factor of the intensity enhancement reached maximum at 8.3. The absolute value of intensity was not calculated yet. The calibration of the detectors and the optical component are still under investigation.

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