Abstract

This article discusses the possibility of coupling a polymerization reaction to an oscillatory kinetic model, complemented by diffusion which can lead to spatial structure. We used three well-known mathematical models of oscillators: a variant of the Rossler multivibrator, a model proposed by Edelstein, and the Oregon Oscillator. One or some of the terms in the equations of these models come from a polymerization reaction, while the other terms of these equations will come from collateral processes. So, almost any reaction could become oscillatory and/or with dissipative structure, adding the adequate collateral processes. The propagation stages are considered as invariants and initiation reactions of order $\alpha = 0$, 1, or 2, and termination reactions of order $\beta = 1$ or 2 are assumed. Except in the case $\alpha = 1$, $\beta = 2$, all six remaining reactions combinations ($\alpha$, $\beta$) can be coupled to least one, and often to several of the models. The effects of destabilization a stable homogeneous steady state by the presence of diffusion is also discussed, which is always be possible.

Keywords: Reaction kinetics, Rossler multivibrator, Oregon oscillator, Oscillatory chemical reactions

Introduction

Theoretical and experimental studies of oscillating reactions have been studied at many laboratories and the interest of this type of phenomena has increased rapidly which led to the discovery of oscillations in biochemical and chemical systems [1]. Interestingly, in the early studies, thermodinamists and applied mathematicians worked out models for such reactions and discussed the possibility of oscillations in homogeneous chemical systems. Later a general criterion was available for the type of chemical reactions that could present undamped oscillations [1].

When the first discoveries in oscillatory chemical reactions [2] and spatial structures were found in chemical systems [3] (e.g., concentration waves), mathematical expressions to fit experimental results were looked for.

Adequate models for the chemical oscillations [4] are provided by non-linear ordinary differential equations when they are considered as precursors of the formal kinetic development. Partial differential equations [5] could explain the appearance of spatial rearrangements when they are applied to the reaction kinetics considering spatial diffusion. There are some kinetic reaction mechanisms, which are in general very complex that can be fit to mathematical models to predict the apparition of oscillations and/or spatial arrangements [6]. Recently, it has been suggested [7] that the mathematical equations of these models could be divided in two groups. Some of these equations will consider the main reactions (object to study), and the rest of them will take in account collateral processes. So, almost any reaction could become oscillatory and/or with dissipative structure by adding the adequate collateral processes. In this paper, examples of the application of this idea to polymerisation reactions, is shown.

General Theory

Polymerisation processes and rate equations [7]

Three steps are considered in the polymerisation reactions:

\[ \alpha X \xrightarrow{k} R \]

where reaction rate is of the form $R \sim X^{\alpha}$, and $a$ can take the values 0, 1 or 2

- **Propagating:** $R + X \xrightarrow{k_p} R$
- **Termination:** $BR \xrightarrow{k_t} \text{products}$

In this case reaction rate is of the form $R \sim X^{\alpha}$ and $\alpha$ can take the values 1 or 2.

By assuming the steady state for the propagators $R$, that is $R = R = \overline{R}$, it is obtained that $R \sim X^{\alpha}$; then the expression for the polymerisation rate $\dot{X}_{pol}$ becomes:

\[ \dot{X}_{pol} = -\alpha k_i X^\alpha - k_j X^\beta \]

where $k_i$ includes the constants $k_p$ and $k_t$.

The term $X$ represents monomer concentration, and the propagation stages are considered as invariants and the different mechanisms of the polymerisations will be denoted using the symbols $[\alpha, \beta]$ [7,8], where the initiation reaction ($\alpha$) can be of zero, first, or second order and the termination reaction ($\beta$) can be of first or second order.

Analysis without diffusion

Ordinary differential equations system [4,9] in two variables, $X$ and $Y$, will be considered

\[ \dot{X} = M(X, Y) \]
\[ \dot{Y} = N(X, Y) \]

General Theory

Chemical Oscillations and Spatial Structures in Polymerisation Reactions

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where the variable $Y$ represents the concentration of a species that appears in a collateral process. The steady states denoted by $SS_{i}$, fulfill the equalities $X = 0, \ Y = 0$, and the $X$ and $Y$ concentrations can be calculated by solving $M(X, Y) = 0, N(X, Y) = 0$ equations. A linearized variational system [10,11] is obtained when considering the perturbations defined as:

$$\begin{align*}
x(t) &= X - X_0 \\
y(t) &= Y - Y_0
\end{align*}$$

the system is of the form

$$\begin{align*}
x &= cX + gY \\
y &= aX + bY
\end{align*}$$

where the coefficients

$$\begin{align*}
c &= \left( \frac{\partial M}{\partial X} \right)_0, \ g &= \left( \frac{\partial M}{\partial Y} \right)_0, \ a &= \left( \frac{\partial N}{\partial X} \right)_0, \ b &= \left( \frac{\partial N}{\partial Y} \right)_0
\end{align*}$$

are the elements of the characteristic no–diffusive matrix

$$M_n = \begin{pmatrix} c - n^2 & g \\ a & b - n^2 \end{pmatrix}$$

This matrix trace is

$$T_n = b + c$$

and its determinant is

$$\Delta_n = bc - a$$

In the neighbourhood of the steady state (SS) the solutions of (3) are in the form

$$\begin{align*}
x &= Ae^{\omega t} \\
y &= Be^{\omega t}
\end{align*}$$

where, $\omega$ are the roots of the equation:

$$\omega^2 - T_n \omega + \Delta_n = 0$$

The steady state is unstable if, at least, one of these roots has a positive real part. It occurs when $\Delta_n < 0$. However, this condition causes the SS to be a “saddle point”, which normally involves explosions, and so, it is discarded as realistic chemical model.

A root with a positive real part can be obtained if $T_n > 0$.

If $T_n = 0$ and the discriminant is:

$$\Delta_n = T_n^2 - 4\Delta_b$$

if $\Delta_n < 0$ then $\Delta_b > 0$, the steady state is a node; if $\Delta_n > 0$ then $\Delta_n < 0$, i the steady state a non–focus.

In all the equations systems that will be examined, the condition $\Delta_n > 0$ will be fulfilled, and the transitions from $T_n < 0$ to $T_n = 0$ and to $T_n > 0$ will be possible, that is, from a stable focus ($T_n < 0, \Delta_n = -4\Delta_b < 0$) to node ($T_n = 0, \Delta_n = -4\Delta_b < 0$) and to unstable focus ($T_n > 0, \Delta_n = -4\Delta_b < 0$). Hopf [12] proved that this bifurcation leads to the apparition of a limit cycle around the unstable focus, with sustained oscillations of the system. In this paper, the Hopf bifurcations existence is discussed and the presence of spatial structures is not probable.

Including the diffusion process

If diffusion effects along one dimension ($r$) are considered, the equations (2) take the form [13,10]

$$\begin{align*}
\frac{\partial X}{\partial t} &= M(X, Y) + D_x \frac{\partial^2 X}{\partial r^2} \\
\frac{\partial Y}{\partial t} &= N(X, Y) + D_y \frac{\partial^2 Y}{\partial r^2}
\end{align*}$$

where $D_x$ and $D_y$ are diffusion coefficients, respectively. In this context, the steady states are named homogeneous steady state (HSS). Their stability can also be studied by linearization, and when suitable perturbations for the desired boundary conditions are defined. In this paper, the Neuman’s conditions [10] (of no–flow in the system boundary) are used:

$$\begin{align*}
x(t, r) &= Ae^{\omega t} \cos nr = X - X_0 \\
y(t, r) &= Be^{\omega t} \cos nr = Y - Y_0
\end{align*}$$

where $n$ is the wave number, which under the conditions $x(0, L) = \pm A$ and $y(0, L) = \pm B$, must satisfy,

$$n = \frac{m\pi}{L}$$

where $L$ is the length of the system and $m = 0, 1, 2, 3, \ldots$ (an integer).

The secular matrix of the linearized system takes the form

$$M_n = \begin{pmatrix} c - n^2 & g \\ a & b - n^2 \end{pmatrix}$$

whose trace is

$$T_n = T_n - (D_x + D_y)n^2$$

and the determinant is

$$\Delta_n = T_n - n^2(D_x + cD_y) + D_xD_yn^4$$

Now, trace and determinant depend on the wave number, $n$. The roots of the secular equation

$$\omega^2 - T_n \omega + \Delta_n = 0$$

will be also dependent on the wave number. As in the no diffusive case, if one of those roots has a positive real part, the homogeneous steady state will be unstable. That can be so because of the conditions $\Delta_n < 0$, or $T_n > 0$. The condition $T_n > 0$ is impossible to be fulfilled, unless in the no diffusive case $T_n > 0$, that is, the homogeneous steady state was already unstable in diffusion absence. So, it is impossible a homogeneous steady state, which fulfilled $T_n < 0, \Delta_n > 0$, to be destabilized by diffusion. Besides, in the case $T_n > 0$, if $\omega$ had a positive real part because of being $T_n > 0$, that positive real part would be $\omega_{Re+} = \sqrt{(T_n + \Delta_n)/2}$ if $\Delta_n \geq 0$; in one or the other case, the higher positive value of this real part will be reached when $n = 0$. That means that in the case of where are several solutions (for some $n$ values), when it is allowed by the limit conditions, the homogeneous solution ($n = 0$), will be amplified faster and predominates above the no homogeneous ones, and the presence of spatial structures is not probable.

The condition $\Delta_n < 0$ has not been considered yet. If $p = n^2$, it can be expressed as a second grade polynomial:

$$\Delta_n = \Delta_n - p(bD_x + cD_y) + D_xD_yp^2 < 0$$

and it is possible that it could be fulfilled in some interval $p < p^*$ (although $T_n > 0$ and $\Delta_n > 0$, case in which it can be said that the homogeneous steady state can be destabilized by diffusion. Independently of $p > p^*$, and $p^* = \Delta_n / (2D_xD_y)$. If is necessary $T_n > 0$ to be firstly fulfilled. Moreover, as $p^* = [(bD_x + cD_y)/(2D_xD_y)] > 0$, it is necessary that $bD_x + cD_y > 0$. Finally, the condition for that interval
to exist is that the discriminant must be positive,
\[ \Delta_{(16)} = (b \cdot D_x + c \cdot D_y)^2 - 4D_x D_y \cdot \text{Det}_0 > 0 \] (17)

If all of those conditions are fulfilled, as \( T_n < T_0 < 0, \text{Det}_n < 0 < \text{Det}_0 \), the inequality \( \Delta_{(16)} > 0 \) will be also fulfilled, and so, \( T_n + \sqrt{\Delta_{(16)}} > T_0 + \sqrt{\text{Det}_0} \).

If the homogeneous steady state is unstabilized for any wave number, \( n \), an initial fluctuation having a Fourier component with that number, \( n \), will force the system to go out from the homogeneous steady state and a new spatial arrangement produced by fluctuations \[14\] will be obtained.

Now, the three mathematical models will be analysed.

Models
Rossler modified model \[7, 15, 16\]

1A. without diffusion:
\[ \begin{align*}
A & \longrightarrow Y \longrightarrow \text{polymer} \\
A + X & \longrightarrow 2X \\
X + Y & \longrightarrow \text{Y + degradation products of X} \\
2X & \longrightarrow D_1 \\
A + D_1 & \longrightarrow 3X \\
2X & \longrightarrow \text{D_2} \\
D_1 + X & \longrightarrow D_3 + \text{degradation products of X} \\
\alpha X & \longrightarrow R \\
R + X & \longrightarrow R \\
\beta R & \longrightarrow \text{polymer}
\end{align*} \]

The kinetic equations:
\[ \begin{align*}
\dot{X} &= A \cdot m_X + a_1 X - m_X Y - 2m_X X^2 + 2m_1 A_1 D_1 - 2m_X X^2 - m_2 X^3 - m_3 D_1 X^2 - k_{RX} X \\
\dot{Y} &= m_X - F_Y \\
\dot{D_1} &= m_X - m_A D_1 \\
\dot{D_2} &= m_X - m_A D_2 \\
\dot{R} &= k_{RX} X - \beta R X
\end{align*} \]

When \( \dot{D}_1 = 0 \), \( \dot{D}_2 = 0 \), \( \dot{R}_0 = 0 \), they are reduced to two equations:
\[ \begin{align*}
\dot{X} &= (-m_1 + m_X) X - m_X Y + X^2 - m_1 X^2 + \mu X^3 - \mu X^4 - m_1 X^2 + k_1 X^n \\
\dot{Y} &= m_X - F_Y
\end{align*} \]

and the following equations are obtained:
\[ \begin{align*}
X &= -k_1 X Y + k_2 X^n - k_3 X^3 + k_4 X + k_5 \\
Y &= k_6 X - k_7 Y
\end{align*} \] (18)

where the values of \( k_1, k_2, k_3, k_4, k_5 \) and \( k_6 \) are defined as follows:

Considering the polymerizations types \( \{0, 1\} \) and \( \{0, 2\} \):
\[ A_i = k_i; A_i = m_1 - k_i = k_i \left( \text{con } k_i = k_i \left( \frac{\mu}{\beta \mu R} \right) \right), \beta = 1, \beta = 2. \]

\[ m_1 = k_1, m_2 = k_2, \frac{m_0 m_1}{m_0} = 1, k_6 = m_1, k_7 = -F \]

for the polymerisation type \( \{1, 1\} \):
\[ A_i = k_i; A_i = m_1 - k_i = m_0 \left( \frac{k_k}{m_k} \right) = m_0 \left( \frac{k_k}{m_k} \right) = k_6, m_1 = k_7, F = k_7 \]

for polymerisation type \( \{2, 1\} \):
\[ A_i = k_i; A_i = m_1 - k_i = m_1 - k_2 = k_1 \left( \frac{\mu}{\beta \mu R} \right) = m_1, m_1 = k_7, F = k_7 \]

\[ \text{when the polymerization type \{2, 2\} is considered:} \]
\[ A_i = k_i; A_i = m_1 - k_i = k_2, m_1 = k_2, m_2 = 2k_1 + k_2 \left( \frac{\mu}{\beta \mu R} \right) = k_1, m_1 = k_7, F = k_7 \]

A. A concrete case \[7\] will be considered. With the following values of the constants: \( k_1 = 1, k_2 = 0.9, k_3 = 1, k_4 = 0.2, k_5 = 0.01, k_6 = 1.5 \mu \) and \( k_7 = \mu \).

The system becomes:
\[ \dot{X} = -XY + 0.9X^2 - X^3 + 0.2X + 0.01 \]
(19)
\[ \dot{Y} = \epsilon (1.5X - Y) \]

This system has an only one homogeneous steady state \( X_0 = 0.271, Y_0 = 0.4065 \) and its secular matrix is
\[ M_0 = \begin{pmatrix} c & = 0.061 \\ g & = -0.7305 \end{pmatrix} \]

with \( T_0 = 0.061 - \epsilon \), will be positive if \( \epsilon < 0.061 \), and \( \text{Det}_0 = 0.3455 \mu > 0 \). So, (19) admits Hopf bifurcation and a limit cycle around the unstable steady state for \( \epsilon < 0.061 \).

1B. Considering that there is a diffusion process, (18) becomes:
\[ \frac{\partial X}{\partial t} = -k_X Y + k_2 X_1 - k_3 X_5 + k_4 X + k_5 + D_1 \frac{\partial^2 X}{\partial x^2} \]
\[ \frac{\partial Y}{\partial t} = k_6 X - k_7 Y + D_2 \frac{\partial^2 Y}{\partial y^2} \]

and using the constants values of Section 1A
\[ \begin{align*}
\dot{X} &= -XY + 0.9X_2 - X_3 + 0.2X + 0.01 + D_1 \frac{\partial^2 X}{\partial x^2} \\
\dot{Y} &= \epsilon (1.5X - Y) + D_2 \frac{\partial^2 Y}{\partial y^2}
\end{align*} \]
(21)

The secular matrix is
\[ M_0 = \begin{pmatrix} 0.061 - D_1 n^2 & -0.271 \\ 1.5 \mu & -\epsilon - D_2 n^2 \end{pmatrix} \]

whose trace is \( T_0 = 0.061 - \epsilon - (D_1 + D_2)n^2 \), and the determinant, \( \text{Det}_0 = 0.3455 \mu - n^2 (0.061D_1 + \epsilon D_2) + D_1 D_2 n^4 \).

The condition \( \text{Det}_0 > 0 \) is fulfilled; the \( D_2 > 0 \), implies that \( 0.061D_1 + \epsilon D_2 \), and (17) requires \( (0.061D_1 + \epsilon D_2)^2 > 4D_1 D_2 \).

For \( \epsilon = 1 \) (value in which \( T_0 < 0 \), \( D_1 = 1, D_2 = 600 \) (adequate units), this condition is also fulfilled. \( p = 0.025 \) could be a possible value of \( p \) that makes \( \text{Det}_0 < 0 \); for that \( p \) value, the real part of \( T_0 \) is \( 0.021/2 \), and its higher than the real part of \( T_{v_p} - (T_{v_p}/2) < 0 \) (since \( \Delta_n < -0.5 < 0 \), the homogeneous steady state was an unstable steady state).

Edelstein model \[7, 17, 18\]

2A. This model is based in an enzymatic mechanism \[18-20\] whose substrate is monomer X. That makes possible to be applied in biochemical polymerisations.
\[ \begin{align*}
A_i & \rightarrow U \\
A_i + U & \rightarrow^{m} 2X \\
X + X & \rightarrow^{n_1} C \\
C & \rightarrow^{n_2} Z + P \\
2X & \rightarrow R \\
R + X & \rightarrow^{n_3} R \\
2R & \rightarrow \text{polymer}
\end{align*} \]

The kinetic equations are:
\[
\begin{align*}
\dot{X} &= m_i XU - 2m_i XY + m_i C + m_i C - k_p R X \\
\dot{U} &= A_i - m_i X U \\
\dot{Z} &= -m_i ZX + m_i C + m_i C - m_i Z Y \\
\dot{C} &= m_i ZX - m_i C + m_i C + m_i Z p \\
\dot{R} &= k_p X^2 - 2k_R R^2
\end{align*}
\]

Which are reduced to two variables if \( Z = 0, R = 0, C = 0 \). Observe that \( Z + C = 0 \), then \( Z + C = B = \text{constant} \), resulting in the following equations:
\[
\begin{align*}
\dot{X} &= k_p X Y - k_2 X^2 - k_4 X + k_4 + D_x \frac{\partial^2 X}{\partial r^2} \\
\dot{Y} &= k_p X Y + D_x \frac{\partial^2 Y}{\partial r^2}
\end{align*}
\]

and with the values for the parameters previously used:
\[
\begin{align*}
\frac{\partial^2 X}{\partial r^2} &= \frac{1}{\mu}(X Y - X^2 - 60(X + 1) + 30) + D_x \frac{\partial^2 X}{\partial r^2} \\
\frac{\partial^2 Y}{\partial r^2} &= 16.858 - X Y + D_x \frac{\partial^2 Y}{\partial r^2}
\end{align*}
\]

The secular matrix is
\[
M_n = \begin{pmatrix}
\frac{0.3474}{\mu} & -n^2 \\
-8.429 & -2 - n^2
\end{pmatrix}
\]

whose trace is
\[
T_n = \frac{0.3474}{\mu} - 2 - (D_x + D_y)n^2
\]

and the determinant
\[
\text{Det}_n = 16.1635 \frac{\mu}{n} - n^2 \left( \frac{0.3474}{\mu} D_y - 2D_x \right) + D_x D_y n^4
\]

\( \text{Det}_n > 0 \) is fulfilled. In order to fulfill the conditions
\[
\frac{0.3474}{\mu} D_y > 2D_x
\]

and
\[
\left( \frac{0.3474}{\mu} D_y - 2D_x \right)^2 > 4D_x D_y 16.1633
\]

the values \( \mu = 1 \) (that leads to \( T_n < 0 \)), \( D_x = 1 \), \( D_y = 600 \), are used. So, the conditions for \( \text{Det}_n < 0 \) are fulfilled if \( p = 0.2 \), value that provides a real part of \( \omega > 0.0181/2 \), higher than the real part of \( \omega_0 \) (without diffusion), \( \omega_0/2 < 0 \), since \( \Delta_0 < 0 \). The homogeneous steady state was a stable focus and it could be destabilized by diffusion.

\textbf{Oregon oscillator [7,21], Stiffly coupled oregonator [22]}

5A. It was conceived first by Field and Noyes [21] as a model developed from the well-known Belousov–Zhabotinsky oscillating reaction [2]. Those authors explained it in three variables, but it is possible to reduce it to two variables assuming that the third one is always maintained in steady state (stiff coupling) [23], without losing its main important characteristics.

Considering the following values [7] for the constants (in adequate units): \( k_1 = 1/\mu, k_2 = 1/\mu, k_4 = 60/\mu, k_6 = 60/\mu, k_8 = 1, k_9 = 30/\mu, k_s = 16.858 and k_p = 1 \), the equations system becomes:
\[
\begin{align*}
X &= 1(XY - X^2 - 60(X + 1) + 30) \\
Y &= 16.858 - XY
\end{align*}
\]

which has a steady state in \( X_0 = 2, Y_0 = 8.429 \). The secular matrix is
\[
M_n = \begin{pmatrix}
\frac{0.3474}{\mu} & \frac{2}{\mu} \\
a & -8.429
\end{pmatrix}
\]

and the trace, \( T_n = (0.3474/\mu) - 2 \), if \( T_n = 0.1737 \), then \( T_n > 0 \). Moreover, \( \text{Det}_n = (16.1633/ T_n) > 0 \). So, if \( \mu < 0.1737 \), a limit cycle around the unstable steady state exits.

2B. If the diffusion is considered, (30) can be written as
\[
\frac{\partial X}{\partial t} = k_p X Y - k_2 X^2 - k_4 X + k_4 + D_x \frac{\partial^2 X}{\partial r^2}
\]

\[ \frac{\partial Y}{\partial t} = k_p X Y + D_x \frac{\partial^2 Y}{\partial r^2} \]

The secular matrix is
\[
M_n = \begin{pmatrix}
\frac{0.3474}{\mu} & -D_y n^2 \\
-8.429 & -2 - D_y n^2
\end{pmatrix}
\]

whose trace is
\[
T_n = \frac{0.3474}{\mu} - 2 - (D_x + D_y)n^2
\]

and the determinant
\[
\text{Det}_n = 16.1635 \frac{\mu}{n} - n^2 \left( \frac{0.3474}{\mu} D_y - 2D_x \right) + D_x D_y n^4
\]

\( \text{Det}_n > 0 \) is fulfilled. In order to fulfill the conditions
\[
\frac{0.3474}{\mu} D_y > 2D_x
\]

and
\[
\left( \frac{0.3474}{\mu} D_y - 2D_x \right)^2 > 4D_x D_y 16.1633
\]

the values \( \mu = 1 \) (that leads to \( T_n < 0 \)), \( D_x = 1 \), \( D_y = 600 \), are used. So, the conditions for \( \text{Det}_n < 0 \) are fulfilled if \( p = 0.2 \), value that provides a real part of \( \omega > 0.0181/2 \), higher than the real part of \( \omega_0 \) (without diffusion), \( \omega_0/2 < 0 \), since \( \Delta_0 < 0 \). The homogeneous steady state was a stable focus and it could be destabilized by diffusion.

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\[
\begin{align*}
\dot{X} &= A_2Z - 2m_2ZX + A_2X - 2m_2X^2 - \alpha k_1X^2 - k_rRX \\
\dot{Y} &= A_2X - m_1Y \\
\dot{Z} &= -A_2Z - m_3Y \\
\dot{\bar{R}} &= k_6X^\alpha - \beta k_1R^\beta 
\end{align*}
\]

Which are reduced to two variables if \( \dot{Z} = 0 \) (stiff coupling), \( \dot{R} = 0 \):

\[
\begin{align*}
\dot{X} &= (A - m_1X)m_Y \\
A + m_2X^2 + A_2X - 2m_2X^2 - \alpha k_1X^2 - k_rR X^{(\alpha/\beta)+1} \\
\dot{Y} &= A_2X - m_3Y 
\end{align*}
\]

Considering polymerizations type \( (0, 1) \) and \( (0, 2) \) and with the definitions:

\[
A_i = k_i, m_i = k_i, m_i / m_i = k_i, k_i = 2m_i + k_i = \Lambda_i, k_i = A_i, k_i = m_i, k_1 = A_1, m_1 = m_1
\]

The kinetic equations take the form:

\[
\begin{align*}
\dot{X} &= Y(k_1 - k_1 X) + k_1 X - k_1 Y \\
Y &= k_1 X - k_1 Y \\
\end{align*}
\]  

Polymerization type \( (1, 1) \) with the definitions:

\[
A_1 = k_1, m_1 = k_1, m_1 / m_1 = k_1, k_1 = 2m_1 + k_1, k_1 = A_1, k_1 = m_1, k_2 = A_2, m_2 = m_2
\]

results in (34).

\[
\begin{align*}
\dot{X} &= Y(k_1 - k_1 X) + k_1 X - k_1 Y \\
Y &= k_1 X - k_1 Y \\
\end{align*}
\]

Polymerization type \( (2, 2) \) with the definitions:

\[
A_2 = k_2, m_2 = k_2, m_2 / m_2 = k_2, k_2 = 2m_2 + k_2, k_2 = A_2, k_2 = m_2
\]

and the determinant, M, is given as

\[
\begin{align*}
\dot{X} &= (k_1 - X) + X - 8.375 \cdot 10^{-6} X^2 + D_1 \frac{\partial^2 X}{\partial t^2} \\
\dot{Y} &= k_1 X - k_1 Y + D_1 \frac{\partial^2 Y}{\partial t^2} \\
\end{align*}
\]

with the following secular matrix

\[
M_n = \begin{pmatrix} 0.9877 - D_n n^2 & -0.9959 \\
\varepsilon & - \varepsilon - D_n n^2 \end{pmatrix}
\]

whose trace is,

\[
T_n = 0.9877 - \varepsilon - (D_x + D_y) n^2
\]

and the determinant,

\[
\text{Det}_n = 0.9877 \cdot 10^{-5} \varepsilon - n^2 (0.9877 D_y - \varepsilon D_x) + D_x D_y n^4
\]

When \( \text{Det}_n > 0 \), the conditions \( 0.9877 D_y > \varepsilon D_x \) and (17) if \( \varepsilon = 1 \) (that means \( T_n < 0 \), \( D_1 = 1, D_2 = 2 \), can be fulfilled; for example, the value \( p = 0.2 \) satisfies the conditions for \( T_n < 0 \). For that value, the real part of \( \omega_n \) is 0.2837/2, higher than the part corresponding to \( n = 0 \), that would be \( T_0/2 < 0 \) (since \( \Delta_0 < 0 \)). The homogeneous steady state is a stable focus.

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