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Smoke on the water: CO and H$_2$O in a circumnuclear disc around a quasar at redshift 2.64

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ABSTRACT

We present an analysis of observations with the Atacama Large (sub-)Millimetre Array (ALMA) of the 380 GHz water emission line and CO (11–10) emission line from MG J0414+0534, a gravitationally lensed dusty star-forming galaxy that hosts a type 1 quasar. We also present observations at 1.6 GHz with global very long baseline interferometry of the radio source. We confirm the previously reported detection of the 380 GHz water line and that the flux density ratio between the two merging lensed images is reversed with respect to the radio/mm continuum. We further find tentative evidence of variability in the integrated line intensity on time-scales of days. We show that the 380 GHz water line has two components of emission: a disc that is around 35 pc in diameter around the quasar, and another component of emission that is offset $\sim$600 pc perpendicular to the disc that lies close to the lensing caustic.

With lens modelling of the multiwavelength data sets, we construct a model for the quasar system consisting of a circumnuclear disc of molecular gas with a size of about 60 pc bisected by radio jets extending to a distance of about 200 pc from the radio core. Our findings suggest that observations with ALMA of high-excitation molecular lines from strongly lensed quasars could allow detailed studies of active galactic nucleus accretion and feedback at the cosmic peak of black hole and galaxy growth.

Key words: gravitational lensing: strong – masers – galaxies: jets – galaxies: nuclei – quasars: general.

1 INTRODUCTION

Water is an abundant molecule in the Universe with a large number of observable transitions from cm to mm wavelengths, whose properties can provide valuable insights into the physical conditions of the interstellar medium (ISM) of galaxies. While H$_2$O emission is weak from the quiescent cold ISM, it is often associated with the dense, warm (or shock-heated) ISM in star-forming regions (e.g. van der Werf et al. 2011; Jarugula et al. 2019) or regions in the vicinity of active galactic nuclei (AGNs; e.g. Moran et al. 1995). Collisional excitation or pumping from a strong radiation field can lead to population inversion, which can create the conditions needed for bright maser emission to occur (Neufeld, Maloney & Conger 1994). Extra-galactic H$_2$O masers can be produced by a large number of individual masers in an intense starburst (Cernicharo, Pardo & Weiss 2006; König et al. 2017). However, they are more typically associated with AGN, where synchrotron emission from the radio core can penetrate the large column densities of dust and gas in the circumnuclear region, and provide sufficient continuum luminosity to reach the critical threshold for maser emission. These are termed ‘megamasers’, having luminosities $>10^6$ times higher than Galactic masers (see Lo 2005, for review).

AGN megamasers have been localized with very long baseline interferometry (VLBI; 6$_{16}$–5$_{23}$ transition; rest-frequency 22.23508 GHz) at cm wavelengths to regions that are typically $<1$ pc from the radio core. These can originate from illuminated gas near or in the accretion disc or torus, where they can be used to constrain the physical environment around the black hole (e.g. Moran et al. 1995; Herrnstein et al. 1998; Klöckner, Baan & Garrett 2003; Kuo et al. 2018). Studies of AGN disc masers have enabled accretion discs to be characterized as geometrically thin discs of dense molecular gas (Neufeld & Maloney 1995; Herrnstein et al. 1998; Klöckner, Baan & Garrett 2003; Kuo et al. 2018). Studies of AGN disc masers have enabled accretion discs to be characterized as geometrically thin discs of dense molecular gas (Neufeld & Maloney 1995; Herrnstein, Greenhill & Moran 1996; Reid et al. 2009; Kuo et al. 2011) with evidence of substructure (Hughes et al. 2008). AGN megamasers can also be produced by interactions between jet plasma and dense clouds of gas, and are therefore a useful probe of jet physics (Peck et al. 2003) and can trace jet-driven outflows of molecular gas (Greenhill et al. 2003). Observations of megamasers at high redshift are, therefore, especially of interest as they can

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be used to constrain the environment around AGN at the peak of cosmic star formation and AGN activity (Madau & Dickinson 2014). 

Thus far, only one 22 GHz H₂O megamaser has been detected at \( z \geq 1 \), from the gravitationally lensed quasar system MG J0414+0534 (\( z = 2.64 \); Impellizzeri et al. 2008; Castangia et al. 2011). The next highest redshift detection is from a type 2 quasar at \( z = 0.66 \) (Barvainis & Antonucci 2005), despite numerous searches of other systems at \( z > 1 \) (Wilner et al. 1999; Livson 2006; Wagg & Momjian 2009; Bennett et al. 2009; McKeen et al. 2011). The megamaser emission from MG J0414+0534 has not been localized, so the origin (whether disc, jet, or starburst) has, thus far, been undetermined. However, the bright synchrotron emission from the radio source and large isotropic line luminosity of \( 10^{19} \) \( \text{L}_\odot \) strongly suggests the megamaser emission is associated with the AGN. The blueshifted, complex velocity profile, initially reported by Impellizzeri et al. (2008), suggested it originates in the radio jet. However, the later tentative detection of a redshifted satellite line by Castangia et al. (2011) may point to a nuclear disc origin.

Variability of 22 GHz H₂O megamasers that are associated with AGN can be caused by fluctuations in the luminosity of the seed photons, the radiative pumping source that causes population inversion (Neufeld et al. 1994; Gallimore et al. 2001), the movement of the masing region (such as in a disc; Moran et al. 1995), or from interstellar scintillation (Greenhill et al. 1997). These variations can occur on time-scales of minutes, in the case of scintillation, to months, in the case of disc rotation. It has also been suggested that microlensing by stars in foreground lensing galaxies can result in significant (extrinsic) variability of lensed 22 GHz H₂O megamasers at high redshift (Gardzen, Lewis & Harvey-Smith 2011). The 15 month monitoring campaign of MG J0414+0534 with the Arecibo telescope that was reported by Castangia et al. (2011) revealed variations in the isotropic line luminosity over 8.2 kpc at the source redshift of 2.64.

Here, we present a re-analysis of the ALMA observations of MG J0414+0534 reported by Kuo et al., and use these in combination with archival 1.6 GHz VLBI data of the radio continuum and ALMA CO (11–10) imaging to investigate the emission properties of the 380 GHz H₂O line. We confirm a robust detection of the 380 GHz H₂O line emission and the reverse in the flux ratio of the two merging lenses images with respect to the other multiwavelength data for this object. We localize the 380 GHz H₂O emission relative to the CO and radio jet components, from which we derive a model for the emission from the system. In Section 2, we describe the observations and data reduction of the ALMA and global VLBI data. In Section 3, we report the velocity structure and variability of the 380 GHz H₂O line, and describe our lens modelling analysis. Finally, in Section 4, we propose a source-plane model to describe the physical origin of the different H₂O components and consider the implications of our results for future searches for mm-wavelength H₂O line emission in the early Universe.

We assume the Planck Collaboration XIII (2016) model of flat cold dark matter cosmology with \( H_0 = 67.8 \text{ km s}^{-1} \text{ Mpc}^{-1}, \Omega_M = 0.31 \), and \( \Omega_{\Lambda} = 0.69 \). For this cosmology, 1 arcsec corresponds to 8.2 kpc at the source redshift of 2.64.

## 2 OBSERVATIONS AND DATA REDUCTION

In this section, we summarize the observations and data reduction of the ALMA and global VLBI data used in this analysis.

### 2.1 1.6 GHz VLBI observations

MG J0414+0534 was observed with the global VLBI array on 2008 June 7, under project code GW019A (PI: Wucknitz). The observations were performed at a central frequency of 1.6 GHz in full circular polarization (RR, LL, RL, LR). The data were correlated with a visibility integration time of 1 s, in eight spectral windows (+/−1217) was performed at intervals of 14 min and standard fringe finder calibrators (3C 286, 3C 48, DA193, and 0528+134) were observed at regular intervals over the 20 h observation. The data were edited, calibrated, and imaged using the Astronomical Image Processing Software (AIPS).

Phase referencing to nearby calibrators (J0412+0438, J0422+0219, and J0409+1217) was performed at intervals of 14 min and standard fringe finder calibrators (3C 286, 3C 48, DA193, and 0528+134) were observed at regular intervals over the 20 h observation. The data were edited, calibrated, and imaged using the Astronomical Image Processing Software (AIPS). The data were manually flagged to remove radio frequency interference and off-source integrations. Observations of the fringe finders were used to remove the instrumental delays. The spectral bandpass was calibrated using observations of DA193 and 0528+134, and the relative amplitudes were set using measured system temperatures and gain curves. Observations of the phase-reference calibration were used to determine the time-dependent phase delays and rates. Several iterations of self-calibration of the target were performed using the line-free spectral windows, down to a solution interval of 60 s, to correct for residual phase and amplitude errors. The self-calibration solutions were applied to the spectral window containing the line emission. The target was imaged and deconvolved in AIPS using a Briggs weighting of the visibilities (shown in Fig. 1). The VLBI components appear structurally similar to those reported by Ros et al. (2000) at 8.4 GHz, and previous imaging of the same data by Volino et al. (2010).

The task IMPFIT within the Common Astronomy Software Applications (CASA) package (McMullin et al. 2007) was used to...
fit 2D Gaussians to the subcomponents of the four lensed images (in the image plane) to extract their positions and uncertainties (listed in Table 1). The radio components are measured relative to the brightest component in image B, for consistency with previous studies (e.g. MacLeod et al. 2013). The relative image positions of the brightest radio component (subcomponent I) are consistent with the relative positions of the optical emission from the quasar (Falco, Lehar & Shapiro 1997), so this is likely emission from (or very close to) the radio core.

2.2 ALMA observations

The ALMA imaging of MG J0414+0534 was taken in two distinct data sets, which we describe separately. The first, taken at 100 GHz, targeted H2O, whereas the second, taken at 340 GHz, targeted high-excitation CO.

2.2.1 100 GHz data

MG J0414+0534 was observed with ALMA in Band 3 under project code 2017.1.00316.S (PI: Kuo). The data were taken in four epochs on 2017 December 13, 15 (two observations), and 18 in an antenna configuration with a longest baseline of 2400 m. The data were correlated with a visibility integration time of 3 s, in four spectral windows of central frequency 92, 94, 104, and 106 GHz, respectively. The 94 and 104 GHz spectral windows covered the redshifted rest frequency of two H2O emission lines (414–321, $v_{\text{rest}} = 380.197 \text{ GHz}$, and 523–616, $v_{\text{rest}} = 336.228 \text{ GHz}$). These spectral windows were each correlated with 1920 channels and 0.94 GHz bandwidth; the remaining spectral windows had 128 channels and 2 GHz bandwidth. J0423−0120 was used as the absolute flux and spectral bandpass calibrator. Phase switching to J0427+0457 was performed to correct the time-dependent phase variations. The total time on-target was 160 min, with approximately 40 min per epoch.

The data were calibrated using the ALMA pipeline within CASA. After confirming the quality of the pipeline calibration, the data were averaged in frequency by a factor of 16 to increase the signal-to-noise ratio per channel and the computational speed during imaging. This resulted in a spectral resolution of 22.6 km s$^{-1}$ channel$^{-1}$. Self-calibration was performed using the continuum spectral windows to spatially align the four epochs and correct for any residual time-dependent phase and amplitude errors. The continuum subtraction was performed by fitting a model to the continuum in the line-free spectral channels and subtracting it from the data in the visibility plane.
The 100 GHz continuum image (not presented here) shows four point sources of emission, consistent with the findings of Kuo et al. (2019) and with the FIR–radio spectral energy distribution that shows synchrotron-dominated emission at this frequency (Stacey et al. 2018). We find no significant variability in the continuum flux density between the four epochs. We do not detect the $^{34}$S–$^{36}$S H$_2$O line, consistent with the results of Kuo et al. (2019).

We confirm the detection of the $^{4}$H$_{26}$S–$^{33}$H$_{24}$ S H$_2$O line; the line intensity and line luminosity (uncorrected for lensing magnification) are given in Table 2. We detect lensed images A1, A2, and B, but do not detect image C (the least magnified image). Moment maps of the line emission from lensed images A1 and A2, using a signal-to-noise ratio cut of 3σ, are shown in Fig. 2.

### 2.2.2 340 GHz data

MG J0414+0534 was observed with ALMA in Band 7 under project code 2013.1.01110.S (PI: Inoue). The spectral coverage included the CO (11–10) emission line ($\nu_{\text{rest}} = 1267.014$ GHz). Details of the observations, data reduction, and imaging are reported by Stacey & McKean (2018) and are not presented here. The continuum imaging reveals four point sources of emission from the quasar and an Einstein ring of thermal dust emission. Imaging of the CO (11–10) shows four compact components at the location of the quasar images. The resolved velocity field from the two merging images (A1 and A2) shows evidence of ordered rotation around ±500 km s$^{-1}$ from the systemic velocity.

The positions of the redshifted and blueshifted components of the line emission were extracted by creating images from selected channels on either side of the systemic velocity. The positions of these components were extracted by fitting 2D Gaussians using the task IMFIT within CASA.

### 3 RESULTS

In this section, we present an analysis of the observed properties of the 380 GHz H$_2$O emission from MG J0414+0534, and the inferred source-plane properties of the various molecular gas components, relative to the quasar and the radio jets.

#### 3.1 Velocity structure

Fig. 3 shows the line profile of the 380 GHz H$_2$O (4$^{14}$S–3$^{13}$S) emission, relative to the CO (11–10) line profile from Stacey & McKean (2018), using an aperture including lensed images A1 and A2, and also from only lensed image A1 in the case of the former. The 380 GHz H$_2$O line profile obtained from only image A1 has a structure suggestive of a double-horn profile, characteristic of emission from a disc. We fit two Gaussian components to the double-horn profile isolated from image A1, which are centred on $-210 \pm 20$ and $490 \pm 20$ km s$^{-1}$, with full width at half-maximum (FWHM) of $630 \pm 50$ and $420 \pm 40$ km s$^{-1}$, respectively. The CO line profile from both A1 and A2 is more peaked around the systemic velocity compared to the total 380 GHz H$_2$O emission, without a double-horn shape. The overall FWHM of the 380 GHz H$_2$O line profile is $\sim 1000$ km s$^{-1}$, comparable to that observed from the CO (11–10) line (i.e. $1100 \pm 100$ km s$^{-1}$), based on a Gaussian fit. The similar terminal velocities suggest that they are associated with the same region of the source.

The double-peaked 380 GHz H$_2$O profile from image A1, under the assumption that the emission is from a disc, suggests a correction to the systemic velocity of $+140 \pm 10$ km s$^{-1}$. The centre of the CO (11–10) profile also suggests a slightly redshifted systemic velocity. We note that the peaks in the 380 GHz H$_2$O line profile are consistent with the blue and redshifted components of the 22 GHz H$_2$O megamaser system reported by Castangia et al. (2011, shown with black crosses in Fig. 3), but with larger associated dispersions.

From the moment maps of the 380 GHz H$_2$O line emission (see Fig. 2), we find that the line emission is only marginally resolved in total intensity for the two merging lensed images (which are also those with the highest magnification; see below). However, there is evidence of resolved structure in the velocity map, where we find a (red) component at around $+100$ km s$^{-1}$, that is, close to the systemic velocity, which is coincident with the peak in the line total intensity, and a spatially offset (blue) component that is at a velocity of around $-300$ km s$^{-1}$. This velocity structure is markedly different from the CO (11–10), where there is evidence of only a single component associated with the AGN. This could be, in part, due to the better angular resolution of the CO (11–10) observations, where the overall velocity structure is less beam-blended. We see that the velocity dispersion of the 380 GHz H$_2$O emission peaks at around 400 km s$^{-1}$ and is co-spatial with the quasar, both of which is consistent with the properties of the CO (11–10) emission. This further suggests that the major component of these two molecular gas tracers are likely associated.

While the angular resolution of the 380 GHz H$_2$O line observations are around a factor of 2 lower than those of the CO (11–10), the water emission cannot be significantly more extended (as it would form an arc or ring), or be spatially offset (which would produce a different image configuration). Assuming a smooth distribution of mass in the lens, these merging images should be almost mirror

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**Table 1.** The radio VLBI and CO velocity components used in our lens modelling. The components are measured by fitting 2D Gaussians to the images using the task IMFIT within CASA. The positions are defined in arcseconds relative to components B–I.

|   | I     | II    | III   | IV    | Red   | CO    | Blue   |
|---|-------|-------|-------|-------|-------|-------|--------|
| A1 $\Delta x$ | $-0.589392 \pm 0.000005$ | $-0.60892 \pm 0.00004$ | $-0.58320 \pm 0.00002$ | $-0.63538 \pm 0.00010$ | $-0.576 \pm 0.003$ | $-0.596 \pm 0.005$ |
| $\Delta y$ | $-1.938372 \pm 0.000009$ | $-1.92669 \pm 0.00008$ | $-1.91770 \pm 0.00005$ | $-1.88383 \pm 0.00017$ | $-1.993 \pm 0.003$ | $-1.863 \pm 0.005$ |
| A2 $\Delta x$ | $-0.723891 \pm 0.000006$ | $-0.73902 \pm 0.00001$ | $-0.69423 \pm 0.00003$ | $-0.74172 \pm 0.000016$ | $-0.739 \pm 0.003$ | $-0.709 \pm 0.005$ |
| $\Delta y$ | $-1.53326 \pm 0.00001$ | $-1.50852 \pm 0.00004$ | $-1.60586 \pm 0.00005$ | $-1.53706 \pm 0.00008$ | $-1.520 \pm 0.003$ | $-1.620 \pm 0.005$ |
| B $\Delta x$ | $\equiv 0.00000 \pm 0.00000$ | $-0.05185 \pm 0.00010$ | $-0.06236 \pm 0.00006$ | $-0.72623 \pm 0.00030$ | $\equiv 0.000 \pm 0.003$ | $\equiv 0.000 \pm 0.005$ |
| $\Delta y$ | $\equiv 0.00000 \pm 0.00003$ | $-0.02566 \pm 0.00007$ | $-0.03189 \pm 0.00010$ | $-0.03282 \pm 0.00014$ | $\equiv 0.000 \pm 0.003$ | $\equiv 0.000 \pm 0.005$ |
| C $\Delta x$ | $+1.3549 \pm 0.00002$ | $+1.3549 \pm 0.00005$ | $+1.34980 \pm 0.00004$ | $+1.35486 \pm 0.000049$ | $+1.345 \pm 0.003$ | $+1.345 \pm 0.005$ |
| $\Delta y$ | $-1.63822 \pm 0.00005$ | $-1.61619 \pm 0.00005$ | $-1.66682 \pm 0.00013$ | $-1.61619 \pm 0.000080$ | $-1.648 \pm 0.003$ | $-1.648 \pm 0.005$ |

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$^1$This non-detection is not surprising as the excitation conditions ($E_B \sim 3000$ K) mean that it is not likely to be detected from an extragalactic source.
images with a similar flux density. Consistent with Kuo et al. (2019), the H$_2$O line emission from image A2 appears around twice as bright as A1: this flux density ratio is reversed compared to the radio/mm continuum, CO (11–10) emission and multiwavelength archival observations (Stacey & McKean 2018). As the 380 GHz H$_2$O emission is spatially coincident with the CO (11–10) and mm-continuum emission from the quasar, small-scale mass structure in the lens is unlikely to be responsible for the different flux ratio as the emission is subject to the same local lens potential. Therefore, the observed reversed flux ratio is more likely due to the source structure rather than the mass distribution in the lens (see below).

We fit Gaussian components to the four lensed images in the 380 GHz H$_2$O line integrated intensity to determine at what level the emission is spatially resolved. Assuming the 380 GHz H$_2$O line emission that is coincident with the quasar images is unresolved and has image flux ratios consistent with the mm-continuum, we fix four Gaussian components to the size and position angle of the synthesized beam, and fix the component positions to the centroids of the mm-continuum emission. We initially allow the flux density of these components to optimize. We then fix the flux density of images A2, B, and C using the flux ratios reported by Stacey & McKean (2018), normalized relative to the optimized flux density of A1. As shown in Fig. 4, the residuals of this model fit have surface brightness features at the 4σ level, suggesting that some 380 GHz H$_2$O line emission is spatially resolved around image A2.

Fig. 5 shows position–velocity diagrams of a slice through images A1–A2 for the 380 GHz H$_2$O line emission, with the contours of the CO (11–10) emission from the same region overlaid. A double-peak structure can be seen within the extent of the CO disc, as suggested by the line profile. The CO emission has a caret shape, due to the fact that the blueshifted emission is closer to the caustic and is more strongly magnified. An additional velocity component from the 380 GHz H$_2$O emission can be seen around image A2 that is more extended than the CO emission and does not have a mirror component in A1. This is the extended emission seen in the moment 0 image and contributes the narrow peak in the line profile at around $-220$ km s$^{-1}$ (see Fig. 3). The anomalous velocity component is not consistent with the caret shape of the CO (11–10) disc and is spatially offset. Therefore, this cannot be emission from the disc.

### 3.2 Variability

As the observations of the 380 GHz H$_2$O line were obtained in four epochs with separations between a few hours and 6 d, we examine the individual observations to look for short time-scale variability. The emission from the merging images (A1 + A2) appears to increase by a factor $1.7 \pm 0.8$ between the observations taken on 2017 December 13 and 15, but this is not significant as the uncertainty is large. We find tentative evidence for variability when comparing the integrated flux from combining the two 2017 December 15 epochs with the combined preceding and succeeding epochs on 2017 December 13 and 18 (combined to match the signal-to-noise ratio of 2017 December 15 epochs). The significance of this variability is at the $2\sigma$ level when we consider the total integrated emission (see Fig. 6, left-hand panel). However, the significance rises to $3\sigma$ when we only consider the integrated flux from image A1 (Fig. 6, right-hand panel). We do not find evidence that the continuum emission is variable over the observations, which is consistent with the results obtained by Stacey & McKean (2018) for the rest-frame 1267 GHz continuum emission from the AGN. We note that, although differences in amplitude scaling could mimic source variability, this would not affect the signal-to-noise ratio of the line emission.

The strongest evidence for variability in the 380 GHz H$_2$O line is seen in the position–velocity diagrams for these observations, which are also shown in Fig. 5. The emission can be seen to be quite different between the data subsets. The emission in the combined data from the epochs on 2017 December 15 is higher overall, but, in particular, the emission from image A1 shows the largest difference. The time delay between the merging images of the quasar is expected to be 0.4 d, based on our lens model (see below). When comparing the two epochs taken on 2017 December 15, the flux ratio appears to switch, with A2 brighter in the first epoch. This accords with our expectation from the lens model (image A2 is the leading image), but these measurements are within the overall uncertainties. Our findings could be interpreted as two components of 380 GHz H$_2$O line emission, one that is variable on time-scales of less than a few days, as seen in image A1, and one that is constant over the time-scale of the observations.

### 3.3 Source-plane structure

In order to determine the structure of the various emission components from MG J0414+0534 in the source plane, we first determine a model for the lens. This was performed using the parametric lens modelling software GRAVLENS (Keeton 2011). We assume a simple model of a singular isothermal ellipsoid with external shear (SIE+$\gamma$) for the primary lens (G1), and a singular isothermal sphere (SIS) for the satellite galaxy (G2). The 1.6 GHz VLBI positions, having the highest precision and largest number of image components, were used to optimize for the best lens model parameters. The flux densities of the components were not used in the modelling, as these are known to deviate from a smooth lens model (e.g. MacLeod et al. 2013). While an improved fit to the image positions may also be obtained with more complex parametric models, such as an ellipsoidal power-law mass distribution (e.g. Spingola et al. 2018), this will not significantly change our interpretation here. The resulting best-fitting parameters are given in Table 3, and are found to be in good agreement with previous modelling of the system (e.g. Ros et al. 2000; MacLeod et al. 2013). As the absolute astrometry between the various data sets is lost after self-calibration, we align them using the compact core emission that is associated with synchrotron emission from the quasar.

### Table 2. Observed line intensities, line luminosities, and continuum flux densities for the ALMA observations, uncorrected for lensing magnification. A 5 per cent error in amplitude calibration should be assumed for the continuum flux density.

| Component | $I_{\text{line}}$ (Jy km s$^{-1}$) | $L_{\text{line}}$ (L$_{\odot}$) | $L_{\text{mm}}$ (K km s$^{-1}$ pc$^2$) | $S_{\text{cont}}$ (mJy) | Reference |
|-----------|-----------------|-----------------|-----------------|-----------------|-----------|
| H$_2$O (414–321) | 1.0 ± 0.1 | (5.3 ± 0.6) × 10$^2$ | (3.0 ± 0.4) × 10$^{10}$ | 27 | This work |
| H$_2$O (616–523) | 0.30 ± 0.03 | (9.5 ± 0.9) × 10$^2$ | (2.7 ± 0.3) × 10$^{12}$ | 71 | Castangia et al. (2011) |
| CO (11–10) | 6.4 ± 0.3 | (1.16 ± 0.06) × 10$^9$ | (1.78 ± 0.09) × 10$^{10}$ | 24 | This work; Stacey & McKean (2018) |
H$_2$O and CO emission from MG J0414+0534

Figure 2. Velocity integrated line intensity (moment 0), velocity field (moment 1), and velocity dispersion (moment 2) of the H$_2$O (4$_{14}$–3$_{21}$) line emission. The moment 1 and 2 images are created by masking emission below $3\sigma$, where $\sigma$ is the mean rms noise per channel. In the moment 0 image, contours showing the continuum are overplotted with peak fractions 0.3, 0.5, 0.7, and 0.9. The moment 1 and 2 images show contours of the moment 0 image at $5\sigma$, $7\sigma$, and $9\sigma$ (where $\sigma$ is the rms noise). The synthesized beam is shown in the lower left corner and has a size of 0.4 arcsec $\times$ 0.6 arcsec.

Figure 3. Line profiles of CO (11–10) (blue) and H$_2$O (4$_{14}$–3$_{21}$). The H$_2$O line integrated over both A1 and A2 is shown in red, and the line from the peak of A1 is shown in green. The black crosses indicate the peak and FWHM of the 22 GHz H$_2$O megamaser components from Castangia et al. (2011). The H$_2$O line profiles are smoothed with a boxcar kernel of width 5 channels (113 km s$^{-1}$) and scaled by a factor of 4, for comparison. The systemic velocity is relative to $z = 2.639$, using the radio definition of velocity.

consistent image separation of these components in each data set supports the assumption that the component we identify as the radio core is coincident with the quasar.

The reconstructed source-plane emission from the 1.6 GHz VLBI continuum and CO (11–10) velocity components are shown in Fig. 7. We find that the jet and counter jet (the components either side of the radio core) form an almost-linear structure that is extended in an east–west direction by around 200 pc in projection. The red and blue CO (11–10) velocity components are almost perpendicular to the radio jets, and form what we interpret as a molecular gas disc with a size that is 60 pc in projection around the quasar. Given the similarity in the position–velocity diagrams of the CO (11–10) and the 380 GHz H$_2$O emission, we assume that the molecular gas disc also contains part of the water emission observed from MG J0414+0534.

Assuming just half of the 380 GHz H$_2$O emission we detect originates in a disc, and assuming the total magnification of the radio core inferred from our lens modelling is representative of the magnification of the H$_2$O disc ($\mu_{\text{tot}} \simeq 40$), we estimate the isotropic line luminosity is $\sim 7 \times 10^5$ L$_\odot$.

To constrain the origin of the anomalous velocity component in the 380 GHz H$_2$O line emission (i.e. emission not from the disc), we place a grid of mock sources on the source plane and compute the positions predicted by our lens model. Fig. 8 shows the image positions produced by the mock sources relative to the radio components. As the anomalous velocity feature is not detected from the emission peak of A1, we exclude all source positions that produce an image below the mid-point between images A1 and A2. We also exclude any positions that are doubly imaged (outside the tangential caustic), as these would not produce an image in the region of interest. Doubly imaged source positions
Figure 4. Gaussian model fit to imaging of the 380 GHz H$_2$O line integrated line intensity. Left to right: the data; the model on the same colour scale as the data; and the residuals (data–model) in units of $\sigma$, where $\sigma$ is the rms noise in the data. The Gaussian model components are fixed to the size of the synthesized beam (PSF) and their positions are fixed to the image positions in the continuum emission. The component flux densities are initially allowed to optimize, but afterwards are fixed to the continuum flux ratios reported by Stacey & McKean (2018) relative to the optimized flux density of image A1. Residuals of this model fit show surface brightness features around image A2 at the 4$\sigma$ level.

Figure 5. Position–velocity diagrams for the 380 GHz H$_2$O line emission from the merging images in the direction A1 to A2 (left to right). The top image shows the data from all four epochs; the middle image is data only from two epochs on 2017 December 15; and the bottom image is data combined from 2017 December 13 and 18. The contours show the CO (11–10) line emission from the same region, smoothed with a Gaussian kernel to match the angular resolution of the 380 GHz H$_2$O line emission. That could produce an image near the merging images would be on the opposite side of the caustics (several kpc from the quasar, as suggested by Kuo et al. 2019), which we do not consider to be feasible. We find that the remaining viable source positions are very close to the lens caustic, which would produce a high magnification of the anomalous velocity component. However, as the position–velocity diagrams (see Fig. 5) and Gaussian source fitting to the images (see Fig. 4) show the anomalous feature is north of A2 and around 0.3–0.4 arcsec from the midpoint, only the green and blue source regions are in agreement with the observations. This suggests that the anomalous 380 GHz H$_2$O feature originates >600 pc east of the quasar and perpendicular to the molecular gas disc.

Assuming that half of the 380 GHz H$_2$O emission originates in the anomalous feature, we take the highest combined magnification of the merging images produced by the mock sources (about 330) to infer a minimum intrinsic isotropic line luminosity of $8 \times 10^4 L_\odot$.

3.4 Size and structure of the molecular gas disc

From our lens modelling, we find that the CO (11–10) originates in a molecular gas disc with a size (diameter) of 60 pc around the quasar. From this, we estimate the enclosed dynamical mass ($M_{\textrm{dyn}}$) of the system using

$$M_{\textrm{dyn}} = \frac{R V_{\textrm{max}}^2}{G},$$

(1)

where $R$ is the 30 pc radius (60 pc diameter) and $V_{\textrm{max}}$ is the maximum rotational velocity. We assume the rotational velocity is the line velocity at the FWHM ($V_{\textrm{obs}} = 550$ km s$^{-1}$), corrected with an inclination angle of 45 deg (i.e. $V_{\textrm{max}} = V_{\textrm{obs}}/\sin 45^\circ$). We find a dynamical mass of $M_{\textrm{dyn}}(<30$ pc) $\sim 3 \times 10^9 M_\odot$, which is only three times larger than the black hole mass determined by Pooley et al. (2007) from a microlensing analysis of the X-ray emission from the quasar accretion disc. This suggests that the kinematics of the CO (11–10) emission are driven by the properties of the central supermassive black hole.

The resolved velocity structure of the CO (11–10) line, presented by Stacey & McKean (2018), shows a radially increasing velocity profile and appears to show a very high central velocity dispersion. However, 2D velocity fields are strongly affected by beam-smearing effects that can lead to an incorrect interpretation of the global dynamical properties (Lelli, Fraternali & Sancisi 2010). Therefore, we defer a more comprehensive investigation of the disc dynamics.
quasar. These observations were first presented by Kuo et al. (2019), a gravitationally lensed dusty star-forming galaxy that hosts a type 1

Figure 6. The normalized, velocity-integrated flux density of the 380 GHz H2O line emission and rest-frame continuum for the four epochs. We show the integrated flux density of both merging images (A1+A2; left) and for image A1 only (right). The 380 GHz H2O line emission is shown in red and the continuum is shown in blue. The orange region shows the combined flux density and uncertainty for both epochs on 2017 December 15; and the green region shows the combined flux and uncertainty for 2017 December 13 and 18.

Table 3. The best parameters of the lens model, based on the positions of the VLB1 image subcomponents. The primary lens, G1, is modelled with an SIE and external shear, and the secondary lens, G2, is modelled with an SIS. b is the Einstein radius of the lens. The positions (x, y) are given relative to the brightest component of image B. e and γ are the ellipticity and external shear, and θe and θγ are their respective position angles east of north.

|       |       |       |    |       |       |
|-------|-------|-------|----|-------|-------|
|       | (arcmin) | (arcmin) | (arcmin) | (deg) | (deg) |
| G1    | 1.076  | +0.482 | −1.279 | 0.235 | −84.16 | 0.098 | 53.85 |
| G2    | 0.196  | +0.913 | +0.150 | −     | −     |      |      |

to a future work that will use higher resolution imaging of the CO (11–10) line.

Assuming the molecular gas disc also contains water emitting regions with a Keplerian rotation and whose enclosed mass is dominated by the black hole (e.g. Kuo et al. 2011), we use the same formalism as equation (1) to estimate the minimum size of the disc. Here, we assume the maximum rotational velocity is half the separation of the line peaks identified in Section 3, corrected for an assumed disc inclination angle of 45 deg. This implies the size (diameter) of the water emitting disc is around 35 pc, that is, about half the size of the CO (11–10) disc. Further high-resolution imaging of the 380 GHz H2O emission with ALMA could confirm that the double-peak structure is indeed part of a disc, as implied by the position–velocity diagrams shown in Fig. 5.

4 DISCUSSION AND CONCLUSION

We have re-analysed spectral line imaging with ALMA of the rest-frame 380 GHz H2O emission from MG J0414+0534 (z = 2.639), a gravitationally lensed dusty star-forming galaxy that hosts a type 1 quasar. These observations were first presented by Kuo et al. (2019), who reported a tentative detection of the line. Kuo et al. found that the broad-line profile made the continuum subtraction challenging, and that the flux density ratio of the line emission in the two merging lensed images was reversed relative to multiwavelength continuum and CO (11–10) observations for this object.

4.1 Origin of the H2O emission

From a careful re-analysis of these data, we find strong evidence that the 380 GHz H2O line detection is robust. We independently verified the calibration and imaging (of both the target and the calibrators), and compared the line profile and position–velocity diagram with those from the CO (11–10) and the 22 GHz H2O megamaser that were previously reported for this object (Stacey & McKean 2018; Castangia et al. 2011). We find that the line width (∼1000 km s⁻¹) and velocity dispersion (∼400 km s⁻¹) are almost identical to the CO (11–10) emission, and there are peaks in the line profile (at −210 and +420 km s⁻¹) that correspond to the velocities of 22 GHz H2O megamaser components. Such similarities provide circumstantial evidence that the detection is robust. However, the strongest evidence comes from the position–velocity diagram, where there is evidence of a resolved velocity structure that is consistent with a rotating disc. This disc has an inferred size of around 35 pc, given the mass of the black hole, which is again similar to the resolved CO (11–10) emission. We also find evidence for an additional 380 GHz H2O component of emission, which, with lens modelling, is estimated to be ∼600 pc from the quasar. The source of the offset feature is unclear; perhaps corresponding to a jet-cloud interaction or a dense region of star formation. Its location close to the lens caustic suggests it could be very highly magnified by a factor of >1000.

Our findings may imply that the host galaxy of MG J0414+0534 contains a significant number of water emitting regions, and our detection of this particular component is due to a chance lensing configuration. A similar scenario was suggested by Kuo et al. (2019) in their analysis of the same data set, however, they proposed three water-emitting regions separated by a larger distance of around 1.5 kpc that are doubly imaged. While the model we propose suggests a chance high magnification for the anomalous component, the doubly imaged regions in the model of Kuo et al. produce much lower magnifications corresponding to higher intrinsic luminosities. Future observations of 380 GHz line emission from dusty star-forming galaxies will determine the rarity of high-luminosity water-emitting regions and test which scenario is more probable.
A2

Figure 7. Lens-plane (left) and source-plane (right) positions of the CO (11–10) and 1.6 GHz VLBI components. We show only the region of the lens plane around the merging images (A1 and A2). The grey lines trace the critical curve on the lens plane and tangential caustics on the source plane. The observed image positions are shown with crosses. The model-predicted positions on the lens and source plane are shown in red and blue circles for the respective CO (11–10) velocity components, and in black circles for the 1.6 GHz VLBI components. The source plane structure suggests a molecular gas disc of CO (11–10) emission bisected by two perpendicular radio jets.

A1

Figure 8. Lens-plane images (left) corresponding to mock positions of the anomalous 380 GHz H$_2$O velocity component on the source plane (right). The radio core and jets are shown with filled and empty stars, respectively, and the grey lines trace the critical curve and tangential caustics. The source-plane positions are shown in kpc relative to the radio core. Sources that produce images below the mid-point of A1 and A2 (shown with the dotted magenta line) are excluded. We also exclude sources that produce fewer than four images, as they would not produce an image on the relevant area of the lens plane or would be a very large distance from the quasar.

4.2 Structure of the circumnuclear gas disc

We extend the investigation of the 380 GHz H$_2$O line emission in MG J0414+0534 by combining these data with the information provided by the CO (11–10) spectral line emission and global VLBI imaging of radio jet components to construct a model for the structure of the lensed quasar. Fig. 9 shows a toy model of the reconstructed source emission (approximately to scale). This model consists of a molecular gas disc around the quasar with 22 and 380 GHz H$_2$O line emission with a size of around 35 pc, with the CO (11–10) gas emission extending to a diameter of 60 pc. Bipolar compact radio jets extend from the quasar radio core to a distance of about 200 pc, almost perpendicular to the gas disc. The
Figure 9. A toy model (approximately to scale) to explain the observed H$_2$O (blue) and CO (11–10) (orange) line emission, relative to the radio core and jets (black) from MG J0414+0534. This model predicts a molecular gas disc around the AGN with a size of 60 pc, with the 22 and 380 GHz H$_2$O emission enclosed in a disc of around 35 pc diameter. An additional component of 380 GHz H$_2$O emission is predicted to form around 600 pc in projection from the quasar, and may be associated with a highly magnified star-forming region that happens to lie close to the lens caustics (grey lines). The radio jets extend out to 200 pc in projection from the radio core, and are not expected to be associated with the highly magnified region of water emission from the system.

other source of 380 GHz H$_2$O line emission is offset by $\sim$600 pc east of the quasar.

We find tentative evidence for variability in the 380 GHz H$_2$O disc at the 3σ level on time-scales of hours to days. Amplitude fluctuations were previously reported for the 22 GHz H$_2$O megamaser in long-term monitoring (Castangia et al. 2011), although the 1 month cadence could mask shorter term variability. A high time variability is inconsistent with our estimate of >35 pc for the diameter of the 380 GHz H$_2$O disc and suggests a much more compact region with a size of just a few light days. Such high time variability has been detected from compact maser discs in the local Universe (e.g. Braatz et al. 2003). Monitoring with a short-term cadence, such as with single dish observations, could determine whether the tentative variability is real.

The compact size of the molecular gas disc and the excitation conditions of the CO (11–10) and water lines ($n_{H_2} \sim 10^6$–$10^8$ cm$^{-3}$; $E_u \sim 300$–$600$ K) suggest that we are probing a dense and warm environment close to the AGN. Around one third of the dynamical mass enclosed by the CO disc is in the central supermassive black hole, perhaps indicating that we resolve a circumnuclear disc or the outer AGN torus. High densities close to AGN are often associated with maser activity in the local Universe. However, we find no clear evidence that the 380 GHz H$_2$O disc observed here originates from a population of masers: the extremely high integrated luminosity is more consistent with observations of thermally excited H$_2$O lines from dusty star-forming galaxies (e.g. Jarugula et al. 2019).

The 380 GHz H$_2$O line is predicted to produce both thermal and non-thermal maser emission, and exhibits a sharp transition between these processes with density (Gray et al. 2016). Further investigations of the 380 GHz H$_2$O line are needed to determine whether it is expected to be associated with thermal emission from dusty star-forming galaxies at high redshift, rather than non-thermal maser emission (or a combination of both).

Circumnuclear discs of dust and molecular gas have been resolved on pc-scales around AGN in the local Universe, which provide insight into the nuclear structure of AGN and the mechanisms of black hole feeding (Gallimore et al. 2016; Imanishi et al. 2018; García-Burillo et al. 2016, 2019; Combes et al. 2019; Impellizzeri et al. 2019). A detection of a compact disc around the quasar of MG J0414+0534 could suggest that the 380 GHz H$_2$O line is a useful tool to investigate the nuclear structure of AGN at high redshift. However, the low surface brightness of this 380 GHz H$_2$O line emission and the chance lensing configuration we propose suggest that spatially resolving the source emission may be difficult. On the other hand, we find that the CO (11–10) emission probes similarly small physical scales and is significantly brighter than the 380 GHz H$_2$O emission. Observations of CO (11–10) (or even higher excitation CO lines), combining the spatial resolution possible with the long baselines of ALMA and a high gravitational lens magnification, will present a promising approach to spatially resolve molecular tori at high redshift. This could allow detailed studies of supermassive black hole accretion and feedback on 10-pc scales at the cosmic peak of black hole and galaxy growth.

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