1. Introduction

Over the last decade, several new low frequency radio telescope arrays have been built where the primary receptor in the array is a dipole-like element. These arrays include the Long Wavelength Array (LWA, Ellingson et al., 2013), The Precision Array for Probing the Epoch of Re-ionization (PAPER, Parsons et al., 2010), the Murchison Widefield Array (MWA, Tingay et al., 2013), and the Low Frequency Array (LOFAR, van Haarlem et al., 2013). These telescopes operate based on phased-array concepts, and use a combination of analog and/or digital techniques to form the primary beam of the telescope array.

A common challenge of phased-array radio telescope systems is the characterization and calibration of the phased-array antenna pattern, which can change significantly depending on the pointing direction on the sky. Substantial effort has been made to form predictive models of the phased-array response (e.g., Bui-Van et al., 2017, 2018; Sokolowski et al., 2017; Sutinjo, O'Sullivan, et al., 2015; Warnick et al., 2018), which in turn requires an accurate electromagnetic model of the elements in the phased-array, including mutual coupling effects. Work to characterize the element or primary beam pattern of the telescope includes the use of Unmanned Aerial Vehicles (UAVs) with radio transmitters, low earth orbit satellites and astronomical sources (Bolli et al., 2016; Jacobs et al., 2017; Line et al., 2018; Neben et al., 2015; Sutinjo, Colegate, et al., 2015).

Having an accurate model of the primary beam, and hence primary beam aperture, is driven in large part by the demanding calibration requirements for studying the faint radio signals from the early universe (e.g., Mesinger, 2019, chap 5). The lessons and techniques learned from the existing radio telescope arrays are also directly applicable to the future Square Kilometre Array Low-frequency array radio telescope (SKA-Low). In this paper, we present the results of using a holographic technique to directly measure the complex aperture illumination pattern for a prototype SKA-Low “station”, which is comprised of 256 antennas working as a phased array, where signals from each antenna are individually digitized and combined digitally. The prototype system used in this paper is described in Section 3, and the specifications of the SKA-Low stations can be found in Dewdney (2016).

2. Holographic Aperture Imaging

Holographic aperture imaging is based on the Fourier relationship between the distribution of the electric field in the receiver’s aperture and its far field electric field (voltage) reception pattern (Thompson et al., 2017, p. 819):
$E_A(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} V_F(l, m) \cdot e^{-j2\pi(lx + my)} \ dldm,$  \hspace{1cm} (1)

where $x$ and $y$ are the east and north coordinates, respectively, in the aperture and measured in wavelength. $l$ and $m$ are the standard radio interferometric direction cosines with their origin at the phase center and are aligned with the celestial Right Ascension and Declination coordinates. $E_A$ is the complex aperture electric field distribution and $V_F$ the far field voltage reception pattern (beam pattern). By taking the two-dimensional Fourier Transformation of the beam pattern, the aperture distribution can be obtained, including magnitude and phase information. The required complex beam pattern can be measured by scanning the beam of the antenna under test and making use of a second receiver, which acts as a reference.

### 2.1. Aperture Imaging for Dish Antennas

The commonly used methodology for dish antennas, which was first described by Napier and Bates (1973), is illustrated in Figure 1. The beam pattern of the antenna under test (AUT) is sampled on a grid of $lm$-coordinates which is centered around an unresolved calibration source. The dish is scanned under the source, so the beam points consecutively toward all grid coordinates during the measurement. At the same time, a second reference antenna stays focused on the calibration source. By cross-correlating the received signals of AUT and reference antenna for every sampling coordinate, a complex quantity is obtained. Because it is proportional to the beam pattern of the AUT, it can be treated as a voltage-like quantity and we can use it in Equation 1 for $V_F(l, m)$. The measurements with dish antennas are usually very sophisticated, because the antennas have to be adjusted mechanically, and the calibration source changes its celestial position during the lengthy process.

### 2.2. Suggested Method for Phased Arrays

Our suggested technique for phased arrays follows the same principle. Similar approaches have been used to characterize time delays between the stations of LOFAR (Salas et al., 2020) and also to calibrate it on station level (Brentjens & Bordenave, 2017). However, the method presented here is a more direct transfer of holography from dish antennas to phased array telescopes. It is assumed that the telescope under test (TUT) digitizes the individual signals of all antennas and captures a short snapshot of them while a strong calibrator source is in the field of view of the telescope. Then, all required information can be gathered from these data, as beams can be formed to all coordinates on the $lm$-grid in post-processing. Furthermore, even without initial phase calibration, the TUT can form an (unprecise) beam toward the calibration source, which can be used as reference beam. Hence, a separate reference antenna is not required (see Figure 2). This aspect of the technique was already proposed in Wijnholds (2017). While we form cross-correlations between beamformed outputs and the reference beam signal, Wijnholds correlates individual antenna signals with the reference beam signal.

For aperture arrays where mutual coupling is minimal, the station beam can be computed using classical phased array theory: the beam is the product of the array factor and the element pattern (identical for all elements). By steering the beam (pointing to different $lm$-coordinates) and
de-embedding the array factor, the element pattern could potentially be established from one calibrator by taking many snapshots separated in time as the source moves through the sky or by using multiple calibrators. However, for arrays where mutual coupling cannot be neglected, the station beam cannot be decomposed into an array factor and element pattern. Instead, it must be computed as the (possibly weighted) sum of the individual embedded element patterns (EEPs) (Warnick et al., 2018). An EEP is the radiation pattern from an individual element in the array, with all the other elements suitably terminated. It is defined to include the phase term resulting from its position in the array, although this may be more conveniently inserted subsequently (although similar to an array factor, these phase terms each multiply a different EEP, rather than the array factor which multiplies the element pattern). To measure the EEPs, all but one antenna must be terminated; as such, there is no way to scan the beam, so for a tightly coupled array, such as SKA-Low prototypes using log-period dipole elements (Davidson et al., 2019), a calibrator in a particular direction provides information on only the EEP in that direction. Either a large number of calibrators or moving calibrators such as UAVs or satellites would be required to establish the full EEPs.

The aperture image, which is the Fourier Transformation of the obtained array factor, contains information about the receive path gains and phases of all antennas and can be used for calibration of the array.

Assume the TUT comprises $N$ antennas, which are positioned at the coordinates $(x_n, y_n)$ around the origin of the array. The coordinates are normalized to the observed wavelength. Antenna $n$ receives the signal $f_n[k]$, where $k \in \{1, \ldots, K\}$ is the sample index. The calibration source is located at the celestial coordinates $(l_{\text{ref}}, m_{\text{ref}})$. By applying the appropriate phase factors $a_n$ to the antenna signals (This process requires prior knowledge of the locations of antennas and mappings between antenna and digital signals, however the technique also has utility in verifying these mappings as discussed in Section 3), a beam is formed toward the calibration source and the reference signal $g_{\text{ref}}[k]$ can be computed:

$$
g_{\text{ref}}[k] = \frac{1}{N} \sum_{n=1}^{N} a_n \cdot f_n[k] \quad a_n = e^{i2\pi(x_n l_{\text{ref}} + y_n m_{\text{ref}})} \quad (2)
$$

The grid coordinates, for which the beam is scanned, are chosen as:

$$
l_p = p \cdot \Delta l \quad p \in \{-M / 2, \ldots, M / 2\} \quad m_q = q \cdot \Delta m \quad q \in \{-M / 2, \ldots, M / 2\} \quad (3)
$$

The spacings in the grid, $\Delta l$ and $\Delta m$, and number of points per dimension, $M + 1$, determine the size and resolution of the aperture image. We discuss the choice of these values below. The scan signals, $s_{pq}[k]$, are calculated for all scanning directions centered around the reference direction:

$$
s_{pq}[k] = \frac{1}{N} \sum_{n=1}^{N} a_n e^{i2\pi(x_n l_p + y_n m_q)} \cdot f_n[k] \quad (4)
$$

Then, a quantity proportional to the sampled array factor of the TUT is computed by cross correlating the reference signal with the scan signal for each grid point. The “Beam-Correlation” is defined as

$$
A[l_p, m_q] = \frac{1}{K} \sum_{k=1}^{K} g_{\text{ref}}[k] \cdot s_{pq}[k] \quad (5)
$$
The two-dimensional Fast Fourier Transformation (FFT) of the Beam-Correlation provides the aperture image.

2.3. Fourier Relationship

A fundamental principle of the discrete Fourier Transformation is that the sampling interval in one domain is reciprocal to the range of covered values in the other domain. In this case, the choice $\Delta l$ and $\Delta m$ determines the area that is covered by the aperture image. The critical angular sampling interval that is necessary to get an image of the whole aperture is set by the size of the aperture and the observation wavelength by the following relationship, where $D_{\text{East}}$ and $D_{\text{North}}$ depict the diameter of the array:

$$\Delta l_{\text{max}} = \frac{\lambda}{D_{\text{East}}} \quad \Delta m_{\text{max}} = \frac{\lambda}{D_{\text{North}}}$$ (6)

According to Equation 3, the number of sampling points in each angular dimension is $M + 1$. With $\Delta l$ and $\Delta m$ being fixed by the size of the aperture, the choice of $M$ determines for what range of angles the array factor is scanned. Also, the number of points in each dimension of the aperture image is $M + 1$. Thus, computing $A[l_p, m_q]$ for a wide range of angles increases the resolution of the aperture image. We note that $A[l_p, m_q]$ can be computed for all $(l_p, m_q)$, including for values >1, not just for combinations that map to the celestial sphere. By sampling the $(l, m)$ grid for values >1, the $(x, y)$ grid resolution can be resolved as required.

3. Measurements on the Engineering Development Array 2

The Engineering Development Array 2 (EDA-2) is located at the Murchison Radio-astronomy Observatory (MRO) in Western Australia. It consists of 256 dual polarization dipole antennas, pseudo-randomly spread on a circular field with a diameter of 35 m. The array operates in the frequency range between 50 and 350 MHz. Groups of 16 antennas are connected to an in-field “SMART-Box” in which the signals are converted to be transmitted via optical fiber to the control building, where they are converted back to electrical signals, amplified, filtered, and digitized (Naldi et al., 2017). Signals from each antenna are passed through an oversampled polyphase filterbank that produces “coarse” channels of $\approx 0.93$ MHz bandwidth (Comoretto et al., 2017). Short dumps of raw voltage data from all antennas for a single coarse channel can be captured and written to disk for offline processing. The EDA-2 has identical station configuration and type of antenna to the EDA-1 (Wayth et al., 2017), with the difference that all antennas are individually digitized in EDA-2.

We applied the technique to approximately 0.5 s of voltage data in a single coarse channel of the EDA-2, captured around noon of 2020-05-20. The data were centered on 149 MHz. The Sun was used as strong calibration source because its angular width is sufficiently small compared to the station beam. The Beam-Correlation is computed for $\Delta l$ and $\Delta m$ chosen according to Equation 6 and $M = 80$. Figure 3 depicts the Beam-Correlation for the data. To reduce ringing artifacts in the aperture image, we multiply the $81 \times 81$-point Beam-Correlation with a two-dimensional Gaussian window (with $\sigma = 1$ in both $l$ and $m$ dimension). Afterward, in order to oversample the aperture image, the Beam Correlation is zero-padded.

Figure 5. Phase plot of the Aperture Image, in degrees. Patches with a corresponding magnitude of less than $-10$ dB are faded. Antenna positions are marked with a circle, the previously flagged antenna with a cross.

Figure 6. The difference in the phase solution of holographic and reference calibration, in degrees. Flagged and low gain antennas are marked in black.
Finally, the Fourier Transform is computed, which is shown in Figures 4 and 5.

In the aperture images shown in Figures 4 and 5, the previously measured antenna positions are marked with a circle. Any antennas that were previously identified as having poor or no signal are tagged with a cross, however, for the sake of illustration, their data have still been included in the process. Malfunctioning antennas would normally be “flagged” and excluded from data processing. In the phase plot, patches with a corresponding magnitude of less than $-10$ dB are faded. The magnitude of the aperture image corresponds very well with the antenna positions. As expected, antennas with poor or no signal appear strongly attenuated or not at all in the magnitude plot.

From the phase distribution (Figure 5), it is possible to assign excitation phases to all antennas, using their known positions in the aperture. If the array is already calibrated, the phase excitation should be equal for all antennas. Without initial phase calibration, the excitation phases from the image provide the receive path phases of all antennas.

To evaluate their validity, they are compared with a reference calibration solution, which was measured by a more conventional method. Visibilities were computed between all antenna pairs then phased toward the Sun. A standard radio astronomy calibration procedure that assumes a compact source at the phase center was used to solve for antenna-based complex gains. Baselines shorter than $5\lambda$ were not used. For each element, the difference between both calibration phases is depicted in Figure 6.

Flagged antennas and antennas with a magnitude of less than $-10$ dB (according to Figure 4) are marked in black. The RMS phase difference between the conventional calibration solutions and the aperture phases was measured to be 9.2°, excluding the flagged antennas.

Figure 7 shows the ratio of amplitude of the calibration gain between the holographic and conventional calibration methods. The overall scaling is arbitrary since the holographic method does not depend on the absolute flux scale. While some systematic variation is present (a small gradient in the gain amplitude across the array), it is clear that the methods produce similar gain amplitude estimates up to a scaling constant.

Figure 8 shows the gain amplitude as a function of antenna index for the two methods. Here the mapping from antenna index to antenna location is not relevant; the point of the figure is that it shows substantial variation in amplitude between antennas. However, the overall pattern is the same for both methods and results in the largely uniform gain ratio shown in Figure 7. We discuss the small gradient in amplitude below.

We do not expect the conventional calibration and holographic methods to produce exactly the same result. With conventional calibration, baselines shorter than $5\lambda$ were discarded so that the solution was not biased by large-scale emission in the sky, whereas the holographic method forms station beams where all antennas are included. A detailed treatment of the errors (especially systematic) is beyond the scope of this paper, however we would expect that since the station beam cannot “discard short baselines” this holographic technique will be more susceptible to systematic errors due to all the other sources (and large-scale emission) in the sky besides the reference source.
To illustrate this point, Figure 9 shows the sky imaged from the data used in this work. The color scale has been deliberately adjusted to highlight weak/faint emission compared to the very bright Sun. Emission from the Galactic plane, Galactic supernova remnants, and some brighter compact sources is apparent.

A related potential bias of the method is due to the autocorrelation of system noise (not signals from the sky) of the reference beam with the steered beam. Because phases are applied to each antenna’s signal to steer the beam, the system noise is effectively re-randomized for each unique pointing direction, and decorrelates on the angular size scale of the station beam. In this situation, the signal in the reference beam is dominated by signals from the sky: the Sun is a $\sim 50 \, \text{kJy}$ source and the array SEFD is approximately 1,800 Jy, however, the method may be biased in the case of a lower signal-to-noise ratio from the reference source.

The bias caused by self-noise and large-scale structure discussed here can of course be eliminated by using a reference beam from a different antenna or station, as is done in conventional dish-based holography. In the context of SKA-Low, this should be possible for stations in close proximity by using an existing calibrated station to generate a reference beam. For this work, the focus is to demonstrate the ability and utility of phased arrays to self-generate a reference beam for the holographic process.

Figure 9. The image of the sky made from the dataset used in this work. The greyscale is in Jy/beam, and has been “burned out” to highlight weak/faint emission. For reference, the Sun (top center) was calibrated as a 51,000 Jy source.

Figure 10. Overlay of an aperture image with an aerial drone photo of the EDA-2. The aperture image is partially transparent, and antenna dipoles can be seen in the locations of the peaks in the aperture as expected.
4. Conclusion

Holographic aperture imaging is mainly associated with dish antennas although the measurement procedure for phased array telescopes with multi-beam capabilities is much simpler. The holographic calibration technique suggested here was successfully tested on a phased array telescope, which is similar in its layout to the envisaged SKA-Low stations.

With less than one second’s worth of voltage data, the technique provides the ability to:

- measure the relative complex gain of each antenna in the aperture array
- verify antenna mappings through the analog and digital signal paths (see Figure 10)
- through regular measurements, monitor the phase or delay of each antenna as a function of time, which can be used to detect changes of delays in the system

Hence, the aperture image provided by this method is a simple tool that can be useful for calibration and maintenance of a phased array telescope.

Finally, the calibration solution from holography takes both the receive path gains and the EEPs for the pointing direction of the calibrator into account. Thereby the receive path gain factor is direction independent, while the contribution of the EEPs depends on the celestial position of the calibration source. Similarly to conventional calibration, the calibration solution from holography incorporates contributions from both EEPs and direction independent receive path gains. By applying the technique with calibration sources at various sky-coordinates and comparing the solutions, it might be possible to extract EEPs for all individual antennas and use them to verify simulated EEPs. However, such exact measurements might require a more complex measurement, such as using a separate antenna to form an independent reference beam.

Data Availability Statement

Code and data used for this project are available at https://doi.org/10.25917/5f53032844ed8.

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