Inability of a graph neural network heuristic to outperform greedy algorithms in solving combinatorial optimization problems like Max-Cut

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MATTERS ARISING from Martin J. A. Schuetz et al. Nature Machine Intelligence
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In Ref. [1], Schuetz et al provide a scheme to employ graph neural networks (GNN) as a heuristic to solve a variety of classical, NP-hard combinatorial optimization problems. It describes how the network is trained on sample instances and the resulting GNN heuristic is evaluated applying widely used techniques to determine its ability to succeed. Clearly, the idea of harnessing the powerful abilities of such networks to "learn" the intricacies of complex, multimodal energy landscapes in such a hands-off approach seems enticing. And based on the observed performance, the heuristic promises to be highly scalable, with a computational cost linear in the input size $n$, although there is likely a significant overhead in the pre-factor due to the GNN itself. However, closer inspection shows that the reported results for this GNN are only minutely better than those for gradient descent and get outperformed by a greedy algorithm, for example, for Max-Cut. The discussion also highlights what I believe are some common misconceptions in the evaluations of heuristics.

Among a variety of QUBO problems Ref. [1] consider in their numerical evaluation of their GNN, I want to focus the discussion here on Max-Cut. As explained in the context of Eq. (7), it is derived from an Ising spin-glass Hamiltonian on a $d$-regular random graph [2] for $d = 3$. (In the physics literature, for historical reason such a graph is often referred to as a Bethe-lattice [3, 4].) Minimizing the energy of the Hamiltonian, $H$, maximizes the cut-size $\text{cut} = -H$. The cut results for the GNN (for both, $d = 3$ and 5) are presented in Fig. 4 of Ref. [1], where they find $\text{cut} \sim \gamma_3 t$ with $\gamma_3 \approx 1.28$ via an asymptotic fit to the GNN data obtained from averaging over randomly generated instances of the problem for a progression of different problem sizes $n$. In Fig. 1(a) here, I have recreated their Fig. 4, based on the value of $\gamma_3$ reported for GNN (blue line). Like in Ref. [1], I have also included what they describe as a rigorous upper bound, $\text{cut}_{ub}$ (black-dashed line), which derives from an exact result obtained when $d = \infty$ [5]. While the GNN results appear impressively close to that upper bound, however, including two other sets of data puts these results in a different perspective. The first set I obtained at significant computational cost ($\sim n^3$) with another heuristic ("extremal optimization", EO) long ago in Ref. [6] (black circles). The second set is achieved by a simple gradient descent (GD, maroon squares). GD sequentially looks at randomly selected (Boolean) variables $x_i$ among those whose flip $(x_i \mapsto \neg x_i)$ will improve the cost function. (Such "unstable" variables are easy to track.) After only $\sim 0.4n$ such flips, typically no further improvements were possible and GD converged; very scalable and fast (done overnight on a laptop, averaging over $10^3 - 10^5$ instances at each $n$, up to $n = 10^5$). Presented in the form of Fig. 1(a), the results all look rather good, although it is already noticeable that results for GD are barely distinguishable from those of the elaborate GNN heuristic.

To discern further details, it is essential to present the data in a form that, at least, eliminates some of its trivial aspects. For example, as Schuetz et al reference themselves, the ratio $\text{cut}/n \sim \gamma$ converges to a stable limit with $\gamma \sim d/4 + P_\ast \sqrt{d}/4 + O(\sqrt{d}) + o(n^0)$ for $n, d \to \infty$ [6], where $P_\ast = 0.7632 \ldots$ [3]. In fact, for better comparison with Refs. [3, 4], we focus on the average ground-state energy density of the Hamiltonian in their Eq. (7) at $n = \infty$, which is related to $\gamma$ via
\[ \langle e_d \rangle / \sqrt{d} = \sqrt{d/4 - \gamma \sqrt{4/d}}. \] 
(The awkward denominator is owed to fact that \( P_d = \lim_{d \to \infty} \langle e_d \rangle / \sqrt{d} \). Also, energy provides a fair reference point to assess relative error because a purely random assignment of variables results in an energy of zero, the ultimate null model. Such a reference point is lacking for the errors quoted in Tab. 1 of Ref. \[3\], for example.)

More revealing then merely dividing by \( n \) is the transformation of the data into an extrapolation plot \[4\] (7): Since we care about the scalability of the algorithm in the asymptotic limit for large problem sizes \( n \to \infty \), which in the form of Fig. \[1\] (a) is out of view, it expedient to visualize the data plotted for an inverse of the problem size (i.e., \( 1/n \) or some power thereof \[8 \[9\]). Independent of the largest sizes \( n \) achieved in the data, it conveniently condenses the asymptotic behavior arbitrarily close to the y-intercept where \( 1/n \to 0 \), albeit it at the cost of sacrificing some data for smaller \( n \). To this end, I propose to plot the data in the finite-size corrections form,

\[ \langle e_3 \rangle_n \sim \langle e_3 \rangle_{n=\infty} + \frac{\text{const}}{n} + \ldots, \quad (n \to \infty). \]  

In Fig. \[1\] (b) we have plotted the same data from Fig. \[1\] (a) according to Eq. \[1\] for \( d = 3 \) (modulo a trivial factor of \( 1/\sqrt{3} \) for better comparison with \( P_+ \)). Stark differences between each set of data appear, since each set converges asymptotically to a stable but distinct limit at \( 1/n = 0 \).

First, we note the addition of a well-known result from replica theory, a one-step replica symmetry-breaking (1-RSB) calculation \[10\] that is expected to yield the actual value for \( \langle e_3 \rangle_{n=\infty} \) (and thus, \( \gamma_3 \)) with a precision of \( 10^{-4} \) (green line), a superior reference value than \(-P_+\) (black-dashed line), valid only at \( d = \infty \) although seemingly sensible in the form of Fig. \[1\] (a). The 1-RSB value is further emphasized by the fact that the EO data (black circles) from Ref. \[4\] smoothly extrapolate to the same limit within statistical errors. Finally, in the form of Fig. \[1\] (b), it becomes apparent that the claimed GNN results (blue line) are systematical far (> 15% at any \( n \)) from optimal (1-RSB, green line) and hardly provide any improvement over pure gradient descent (GD, maroon squares). It appears that the GNN learns what is indeed the most typical about the energy landscape: the vast prevalence of high-energy, poor-quality metastable solutions that gradient descent gets trapped in, missing the faint signature of exceedingly rare low-energy minima. In fact, extending GD by a subsequent 5\( n \) spin flips, say, each flip adjusting one among the least-stable spins (even if not always unstable), allows this greedy local search to explore several local minima, still at linear cost. The results of that simple algorithm, also shown in Fig. \[1\] (b) (diamonds), already reduce the error to \( \approx 6\% \) across all sizes \( n \), a considerable improvement on the GNN results in Ref. \[1\] and still better than an improved version, GraphSAGE, the authors mention in their response (orange line).

In conclusion, the study in Ref. \[1\] exemplifies a number of common shortcomings found in the analysis of optimization heuristics (see also Ref. \[7\]): (1) Reliance on rigorous but rather poor and often meaningless bounds, as provided by the Goemans-Williamson algorithm in this case, instead of using the much more relevant results (albeit as-of-yet unproven) from statistical physics, (2) using an obscure presentation of the data, (3) lack of state-of-the-art comparisons across different areas in science, and (4) lack of benchmarking against trivial, baseline models such as gradient descent or greedy search we presented here. On such closer inspection, the proposed GNN heuristic does not provide much algorithmic advantage over that base line. It is likely that these conclusions are not isolated to this specific example but would also hold for Max-Cut at \( d = 5 \) and for the other QUBO applications discussed in Ref. \[1\], as the concurrent comment by Angelini and Ricci-Tersenghi (arXiv:2206.13211) indicates.

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[2] Technically, their Hamiltonian in Eq. (7) pertains to an antiferromagnet instead of a spin glass, but on such random graphs, both are equivalent \[11\].
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