Clay Swelling: Role of Cations in Stabilizing/Destabilizing Mechanisms

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ABSTRACT: The stepwise hydration of clay minerals has been observed repeatedly in studies, but the underlying mechanism remains unclear. Previous numerical studies confirmed the presence of one-water layer (1W) and two-water layer (2W) hydration states. However, the undisturbed transition between these hydration states has never been captured. Using molecular dynamics simulation, this study (i) simulated for the first time the free 1W→2W transition during clay hydration and (ii) identified the underlying mechanism to be the detachment of cations from the clay surface and the formation of a shell of water molecules around the cation. The swelling dynamics of clay was found to be affected by the clay charge, mineralogy, and counterions through complex cation–clay interactions, cation hydration capacity, and cation migration rate.

INTRODUCTION

Clays, such as montmorillonite (MMT) and beidellite (BD), are important constituents of soil and have wide engineering applications, including building material, environmental remediation, catalysis, and geological disposal of nuclear waste.1−4 MMT and BD carry negative charges in their platelike layers due to the isomorphous substitution of silicon (Si) or aluminum (Al) atoms with atoms of a lower valence. To balance out the charge, these substitutions are associated with positively charged cations within the interlayer space.5 In the presence of water, the interlayer cations attract water molecules, pushing clay layers apart, in a process generally described as clay hydration or clay swelling.6,7 However, questions remain on the molecular dynamics driving clay swelling, particularly related to the molecular interactions that affect the movement of cation and water molecules in the interparticle layer of expansive clays.

Clay swelling occurs in two regimes, crystalline and osmotic, depending on the d-space of the clay particle. The d-space is defined for clay particles as the sum of the interlayer distance and the thickness of one clay layer. Crystalline swelling occurs when the d-space increases from 10 to 22 Å, and osmotic swelling occurs when the d-space swells beyond 22 Å.8 Crystalline swelling has been observed to occur in a stepwise fashion.9 The d-space changes abruptly from dry to a one-water layer (1W) hydration state, as well as from a one-water layer to a two-water layer (2W) hydration state.8−10 Numerical simulations have been performed to investigate aspects of the hydration process of clay.9−18 However, previous studies have either (i) inserted water molecules into the space confined by clay layers infinitely extending in horizontal directions,7,16,19−21 (ii) steered molecular dynamics by imposing extra expanding force or potential,15,18 or (iii) focused on swelling dynamics at a larger d-space (i.e., 18.8 Å in ref 11). There remain significant gaps in our scientific understanding of (i) “free” swell from the 1W to 2W hydration states, in which water molecules are allowed to move freely from bulk water to the interlayer space, (ii) the underlying molecular mechanism of stepwise swelling, and (iii) the influence of clay charge, mineralogy, and interlayer cations on

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the “free” swelling of clays. To address these gaps, molecular dynamics simulation on the swelling of clay particles was performed in this study.

**MATERIALS AND METHODS**

Recent development in the CLAYFF force field allows clay particles in a molecular dynamics simulation large-scale atomic/molecular massively parallel simulator (LAMMPS) to have protonated edges. With CLAYFF, water molecules can move freely in and out of the clay particle, something that was not possible in earlier force fields. In this study, the mechanisms underlying stepwise swell were modeled as a tetrahedral (Figure S1 in the Supporting Information). Each layer was investigated using a two-layer model placed in a water box. The top layer was fixed throughout the simulation. The bottom layer was kept fixed throughout the simulation. The top layer was fixed for the first 1 ns of simulation to allow the system to reach equilibrium and then was allowed to move freely vertically, in which case the clay swells naturally. No other restriction was applied to the system (preparation of clay sample and model system are provided in the Supporting Information).

**RESULTS AND DISCUSSION**

Relaxing from an initial d-space of 11.9 Å, stable hydration states and stepwise swell were successfully reproduced for MMT (Figure 1). Stable 1W states were observed for Na with two and four counterions in the interlayer space (Figure 1a,b) and for K with 12 counterions (Figure 1g). Stepwise swell from the 1W to the 2W state was observed for Na with 10 and 12 counterions (Figure 1e,f), and with six counterions for Ca (Figure 1h). There was no evidence of stable, higher hydration states. In simulations relaxed from a higher initial d-space, i.e., at 15.8 and 18.8 Å, Ca6 reached a stable 2W state for both initial d-spaces, and Na12 reached a stable 2W state at a 15.8 Å initial d-space (Figures S4 and S5 in the Supporting Information). The 1W hydration states of Na8 and Ca6 were very short, indicating a direct swell to the 2W hydration state. Continued osmotic swelling was observed for Na with 6, 8, and 10 counterions. The 1W–2W swell occurs rapidly, i.e., in less than 0.2 ns, leading to a sudden change in the measured d-space, which is distinct from the continuous change in the d-space in osmotic swelling (i.e., in Na6, Na8, and Na10). The d-space of 1W and 2W hydration states agree well with the experimental observation and previous numerical studies (Figure 1). In particular, the 1W and 2W d-space of Na12 aligns closely with prior molecular simulation studies (Figure 1f), i.e., 12.57 Å for 1W and 15.03 Å for 2W. The simulations reproduce previously observed swelling behaviors of Na-MMT, specifically the expansion to 2W and the higher d-space for Na-MMT of medium and high charges (Figure 1c–f).

The simulations indicate that the mechanism of crystalline swelling is the detachment and hydration of the counterions. The coordination of water oxygen (Ow) with counterions was observed to begin before the abrupt increase in the d-space and continued during the swelling process, along with a decrease in the number of coordinated clay oxygen (Ob, Figure 2b). Ow and Ob within a 3.0 Å zone from the counterions were correlated between these data sets (Figure 3). A shell of six Ow at the 2W state. The same correlation was found for the 1W to the 2W state was observed for Na, with 10 and 12 counterions. The 1W hydration states of Na8 and Ca6 (not presented here), which supports previous studies reporting the influence of interlayer cation hydration on the crystalline swelling of clay.

While the hydration of counterions is the mechanism generating crystalline swelling, the counterions are still responsible for the dominant attractive force stabilizing the 1W and 2W hydration states. An analysis of the force imposed on the top clay layer identifies an attractive force (F_{1W} < 0) from the interlayer cations, which has a greater magnitude than the repulsive force from the bottom layer (F_{2W} > 0 and |F_{2W}| > |F_{1W} |) (Figure 4a,b). A third force stems from the interaction between the top clay layer and the surrounding water molecules, which can be either repulsive or attractive depending on their relative location. The water molecules within the interlayer space, however, always provide a strong repulsive force on the top clay layer (F_{2W} > 0, Figure 4a,b). In the process of 1W–2W swell, water molecules are attracted to...
coordinate with the interlayer cations. As a result, $F_{TW}$ continuously increases and pushes up the top clay layer (Figure S2 in the Supporting Information). The decomposition of these forces shows that in the first two forces, the Coulombic interaction dominates over pairwise interactions, whereas in the third force, only pairwise interactions are involved (Figure S3 in the Supporting Information). The swelling state of clay particles depends on the competition between these three contributions.

Clay with a higher negative charge (i.e., a greater number of substitutions) experiences a stronger repulsive/attractive force from the bottom clay/cation, but it also experiences a stronger repulsive force from the hydration of interlayer Na (Figure 4a). The full hydration of Na attracted six O$_w$ to form a shell surrounding one single Na (Figure 3). The strong repulsive force exerted on the top layer in this process is indicated by the number of missing coordinated O$_w$. At the 1W hydration state, the repulsive force increases for the clay of a higher negative charge (Figure 4c) as more interlayer Na attracts water and becomes fully hydrated. As a result, the equilibrium state of the Na$_6$, Na$_8$, Na$_{10}$, and Na$_{12}$ swelled beyond the 1W hydration state. A smaller 1W $d$-space was observed for the clay of a higher negative charge (Figure 1), which agrees with the experimental observation and is due to a stronger clay−cation attraction.35 At the 2W state, the stronger clay−cation interaction in the Na$_{12}$ constrained a 2W state (Figure S), whereas Na$_6$, Na$_8$, and Na$_{10}$ continued to swell. For Na$_{10}$, the diffusion of Na from the interlayer region to bulk water leads to weaker constraining forces from the cation to both top and bottom clay layers (Figure 5a−c), which, in turn, reduces the resistance to the migration of water molecules into the interlayer region (i.e., driven by the osmotic pressure) and the osmotic swelling of clay. Alternatively, the stronger force and interaction between clay and cation in Na$_{12}$ restrains Na from moving out of the interlayer region and prevents osmotic swelling, as suggested by previous experimental studies.26 The various hydration states of medium-charged clay and high-charged clay result from the competition between osmotic pressure and cation−clay interaction, which requires further molecular dynamic studies to verify.

In addition to the hydration capacity of the cation, the swelling behavior of Na-, K-, and Ca-MMT varies due to the mobility of the counterion in the interlayer space. The number of coordinated O$_w$ attracted to the cations at its 1W state differs between Ca, Na, and K (Figure 4d), indicating a decreasing hydration capacity.17 This behavior is supported by the radial distribution function of cations and O$_w$ (Figure 4g). Ca had the greatest density of coordinated O$_w$, with K having the lowest and Na having an intermediate density of coordinated O$_w$. The migration of cations (depending on the type of cation) within the interlayer space is greatly inhibited by the clay (Figure 4e,f). Among these cations, the inhibition effect is strongest on K, as its migration rate drops most from the bulk water environment to the interlayer space.36 As a combined result of weak hydration capacity and strong inhibition on migration, K-MMT had the most stable 1W hydration state (Figure 1g).33 The strong hydration capacity of Ca led to a quick swell of Ca-MMT to the 2W state (Figure 1h), whereas its low migration rate prevented it from swelling to higher $d$-space (Figure S4 and 5c).39 The intermediate hydration capacity of Na led to a slow 1W$\rightarrow$2W transition of Na$_{12}$ (Figure 1f), but at a higher $d$-space, the stabilizing effect of Na was weakened by its fast migration rate (i.e., osmotic swelling of Na$_{12}$ when relaxed from 18.8 Å).

Figure 2. Detailed 1W$\rightarrow$2W swell process of Na$_{12}$ (the magenta rectangle in Figure 1f): (a) the evolution of $d$-space and (b) the evolution of coordinated O$_w$ (black curve) and O$_b$ (blue curve) number. Distribution of O$_w$ (red curve) and Na (blue curve) within the interlayer space at the $d$-space of (c) 11.9 Å, (d) 12.57 Å (1W), and (e) 15.03 Å (2W) are present. The slightly lower peaks at a higher $z$ position in d and e are due to the fluctuation of the top layer.

Figure 3. Correlation between the $d$-space and the average number of coordinated O$_w$ (a) and O$_b$ (b) for each cation for MMT.
The hydration of Na-BD relaxed from 11.9 Å $d$-space and reached stable 1W states for all charges (Figure 6a–d). The 1W $d$-spaces of BD (i.e., 13.04 and 13.1 Å for BD-Na4 and BD-Na12, respectively, Figure 6) were larger than those of MMT (i.e., 12.75 and 12.57 Å for Na4 and Na12, respectively, Figure 1). In contrast, the 2W $d$-spaces of BD generated with a higher initial $d$-space were smaller, i.e., around 14.3 Å for BD (Figures S4 and 5d,e) compared with more than 15.0 Å for MMT (Figures 1c–f, S4, and 5a–c). This reverse in the trend has also been observed by a previous modeling study and supported by experimental work. Compared to MMT, $F_{TB}$ and $F_{TC}$ were stronger in BD and the interlayer Na migration

Figure 4. Force analysis showing (a) forces on the top clay layer for Na-MMT with various clay charges (black crosses stand for $F_{TW}$ at the 2W state, and negative/positive value indicates an attractive/repulsive force on the top layer); (b) forces on the top layer for Ca-, Na-, and K-MMT with 6, 12, and 12 counterions, respectively; (c) total coordinated $O_w$ for interlayer Na (black dashed curve indicates the missing number of $O_w$ to be attracted if all Na are fully hydrated); (d) total (black) and mean (red) number of coordinated $O_w$ for Ca, Na, and K extracted at the 1W hydration state; (e) migration rates of interlayer Na, with red, blue, and black color distinguishing cation in the interlayer region, 1 nm away from clay and bulk water; (f) migration rates for interlayer Na, Ca, and K; (g) the radial distribution function of cation and $O_w$. All results presented were extracted from 1W hydration states unless specified.

Figure 5. Comparison of swelling for Na10 (a–c) and Na12 (d–f) showing (a) the interaction energy of clay and interlayer cation (black) and swelling curve (gray); (b) forces from interlayer cation onto the top clay layer ($F_{TC}$, red) and the bottom clay layer ($F_{BC}$, blue); (c) number of interlayer $O_w$ (black) and Na (blue).
rate was lower (Figure 6f), which is due to the shift from the central negative charge in MMT (as O\(^{-}\)Al replaced by Mg) to a surface-negative charge in BD (as T\(^{-}\)Si replaced by Al). The interlayer Na density profile of BD was flat and loosely split into two peaks, in contrast to a single peak profile in MMT, indicating that Na in BD is more attracted to the clay surface (Figure 6h). BD with 12 K and six Ca as counterions also stabilizes at 1W states (Figure S4).

In this study, free swell was investigated from the 1W hydration state, not 0W, because of the presence of an energy barrier for the 0W–1W swell.\(^{12,18}\) At the 1W state and a higher \(d\)-space, the protonated clay edge permits the diffusion of water molecules and cations, but, at the 0W state, exposed \(-OH\) groups form hydrogen bonds in between adjacent layers. According to Ho et al.,\(^{12}\) these hydrogen bonds act like a gate preventing the inflow of water molecules into the interlayer space and impose an energy barrier for 0W–1W swell. Consequently, simulation with the current CLAYFF is not able to simulate the 0W–1W swell. Fixing the bottom layer and controlling the direction of the swell have been used to estimate the osmotic swelling pressure of clay\(^{13,27}\) and are first used to investigate the free 1W–2W swell in this study. It is a stepwise improvement in the modeling. Future work should investigate the chemistry and interactions of the surrounding water with the counterion and its effect on clay swelling.

**CONCLUSIONS**

Using a molecular dynamics approach, this study simulated the free expansion of clay layers during hydration and identified the key mechanisms for the stepwise swell from the 1W to the 2W hydration state. Crystalline swelling is triggered by the detachment of cations from the clay surface and the formation of a shell of water molecules around the cations in the interlayer space. The full hydration of cations causes the clay to swell from the 1W to the 2W hydration state, but it also generates the attractive force that stabilizes both states through clay–cation Coulombic attraction. Importantly, the swelling dynamics of clay is affected by the clay charge, clay mineralogy, and counterions through complex cation–clay interactions, cation hydration capacity, and cation migration rate. Future work should investigate the chemistry and interactions of the surrounding water with the counterion and its effect on clay swelling.

**ASSOCIATED CONTENT**

* Supporting Information
The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsomega.1c04384.

Full computation methods and details on the hydration of Montmorillonite and Beidellite (PDF)

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Notes

The authors declare no competing financial interest.

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