Fast variability as a tracer of accretion regimes in black hole transients

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ABSTRACT
We present the root mean square (rms)–intensity diagram for black hole transients. Using observations taken with the Rossi X-ray Timing Explorer, we study the relation between the rms amplitude of the variability and the net count rate during the 2002, 2004 and 2007 outbursts of the black hole X-ray binary GX 339–4. We find that the rms–flux relation previously observed during the hard state in X-ray binaries does not hold for the other states, when different relations apply. These relations can be used as a good tracer of the different accretion regimes. We identify the hard, soft and intermediate states in the rms–intensity diagram. Transitions between the different states are seen to produce marked changes in the rms–flux relation. We find that one single component is required to explain the ~40 per cent variability observed at low count rates, whereas no or very low variability is associated to the accretion-disc thermal component.

Key words: accretion, accretion discs – black hole physics – binaries: close – X-rays: binaries – X-rays: individual: GX 339–4.

1 INTRODUCTION
Black hole X-ray transients (BHTs) are observed in different states during their outburst evolution. They correspond to particular spectral and timing properties seen in the observations. The physical interpretation of the states is still strongly discussed, but they are probably associated with different accretion regimes. A hard state [historically known as the low/hard state (LHS)] is generally observed at the beginning, and at the end of the outburst. It is roughly characterized by a power-law-shaped energy spectrum with a photon index of ~1.6 (2–20 keV) and a high level of aperiodic variability [root mean square (rms) amplitude above ~30 per cent]. In the middle of the outburst, the energy spectrum is dominated by a soft thermal component and almost no variability is seen. During this soft state [historically known as the high/soft state (HSS)], a hard tail up to ~1 MeV is also present (Grove et al. 1998). In contrast, a high-energy cut-off (≤200 keV) is observed in the LHS (e.g. Wilms et al. 2006; Motta, Belloni & Homan 2009, hereinafter M09). From a physical point of view, the thermal component present in the HSS is usually associated with emission from an optically thick accretion disc, whereas the emission in the hard state is believed to arise from a ‘corona’ of hot electrons, where seed photons from an optically thin accretion disc are up-Comptonized. Synchrotron emission from a jet, which is known to dominate the radio and infrared spectrum during the LHS (e.g. Fender 2006; Russell et al. 2006), could also have a significant contribution to the hard X-ray emission (Russell et al. 2010).

In between these two ‘canonical’ states, the situation is rather more complex. Hard-to-soft and soft-to-hard transitions are observed in relatively short time-scales (hours/days) as compared to the ones seen for the canonical states (weeks/months). During these transitions, both timing and spectral properties change dramatically, leading to different state classifications (see e.g. Belloni 2010 for a review). Homan & Belloni (2005) and Belloni et al. (2005) (hereinafter B05) identified two additional states, the hard-intermediate state (HIMS) and the soft-intermediate state (SIMS), based on spectral and timing properties [see e.g. Casella et al. 2004 for different types of quasi-periodic oscillations (QPOs)]. McClintock & Remillard (2006) proposed a more quantitative but model-dependent classification, which requires a steep power-law state and an intermediate state in addition to the two canonical states (cf. M09).

A fundamental tool for studying the evolution of BHTs is the hardness–intensity diagram (HID), where the evolution of the spectral hardness as a function of the acretion rate can be followed (see e.g. Homan et al. 2001). The hard–soft–hard usual evolution and the main transitions become apparent in the HID. The latter are usually not associated with major changes in flux, with the final soft-to-hard transition occurring at count rates approximately one order of magnitude lower. While the HID provides a general description of the BHT evolution, it is not enough for detailed studies (e.g. it is not able to establish with accuracy when the main transitions occur). Complementary tools, like power density spectra (PDS), the hardness–rms correlation (e.g. B05) and multiwavelength

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observations, are used for a more complete description of the different states and transitions. In this paper, we present a new tool, the rms–intensity diagram (RID), which allows one to map BHT states by looking only at the evolution of the count rate and rms (i.e. without spectral information). As a first step we have studied the case of the BHT GX 339−4, in which detailed observing campaigns have been performed during the multiple outbursts observed with the Rossi X-ray Timing Explorer (RXTE).

GX 339−4 is a low-mass X-ray binary harbouring a \( \gtrsim 6 \, M_\odot \) accreting black hole (Hynes et al. 2003; Muñoz-Darias, Casares & Martínez-Pais 2008). Since its discovery (Markert et al. 1973), the system has undergone several outbursts, becoming one of the most-studied X-ray transient. The intense monitoring carried out by the RXTE during the 2002, 2004 and 2007 outbursts has yielded detailed information on the evolution of black hole states along the outbursts (see e.g. B05). Here, we use this rich data set to study the evolution of the flux and variability along the outbursts.

2 OBSERVATIONS

We have used all the RXTE public archival observations from the 2002 (206), 2004 (295) and 2007 (239) outbursts of GX 339−4. For the analysis, we have considered only the data from the Proportional Counter Array (PCA). ‘Goodxenon’, ‘event’ and ‘single-bit’ data modes are used for the variability. Count rates have been computed using PCA Standard 2 mode data corresponding to the PCU unit #2. The analysis has been performed using HEASOFT v6.7 and a custom timing software running under IDL.

3 THE RMS–INTENSITY DIAGRAM FOR THE 2007 OUTBURST

The RID for the 2007 outburst of GX 339−4 is presented in Fig. 1 as filled dots. For each observation, the net count rate corresponds to PCA channels 0–35 (2–15 keV). Only entire RXTE observations are initially considered, i.e. every dot corresponds to one observation. PDS for each observation have been also computed using the procedure outlined in Belloni et al. (2006). We have used stretches 16 s long and the same energy selection as for the count rate. The total power was computed within the frequency band 0.1–64 Hz and the fractional rms was calculated following Belloni & Hasinger (1990). The absolute rms displayed in the diagram (Fig. 1) is obtained by multiplying fractional rms by net count rate (PCU #2) for each observation. The black solid line joins observations contiguous in time starting from the observation marked with a cross (i.e. observation #1). We find that the source describes a continuous hysteresis pattern in the anticlockwise direction. Four different regions can be distinguished in this diagram.

Figure 1. rms–intensity diagram for the 2002, 2004 and 2007 outbursts of GX 339−4. The dotted lines represent the 1, 5, 10, 20, 30 and 40 per cent fractional rms levels. The grey area marks the part of the diagram where solely type-B QPOs (Casella et al. 2004) are seen. The dashed line joins two consecutive observations separated by 33 d due to observability constraints. The cross corresponds to the first observation of the 2007 outburst.
3.1 The hard line

Starting from the first observation (cross in Fig. 1) and following the solid line, we see the rms increasing with the count rate. This linear trend is seen during the first 18 observations (i.e. \(\sim 40\) d). Following the state classification reported by M09, all these observations belong to the canonical LHS. In the context of this diagram, we will call this linear relation ‘hard line’ (HL). A continuous increase in hardness (M09) is observed during this HL. Weak, type-C QPOs appear in the PDS of the observations corresponding to the upper part of the HL without modifying the rms–flux relation. They become strong once the system shows fractional rms of \(\sim 30\) per cent, and it is close to abandon the HL. In Fig. 2, we plot the RID but using fractional instead of absolute rms. In this representation, the HL we see in Fig. 1 corresponds to the first 18 points starting from the cross, where the fractional rms drops from 40 to 30 per cent.

3.2 State transitions

From observation #19 to #23 (\(\sim 3\) d), the rms starts to decrease, whereas the flux is still increasing. Therefore, the system abandons the HL and follows a different linear relation. From observation #24 to #29 (\(\sim 3\) d), a sudden increase in the count rate is observed, whereas rms is almost constant within \(\sim 4\) count s\(^{-1}\). The spectral fits performed by M09 require a thermal blackbody component from observation #25 (see Fig. 1). Thus, we associate this increase in flux at constant rms to the appearance (2–15 keV band) of an optically thick accretion disc with a very low variability level.

Until observation #32, type-C QPOs are observed in the corresponding PDS. From observation #32 to #33, the rms decreases dramatically, crossing the 10 per cent fractional rms line. A type-B QPO is observed in the PDS. Although the change in rms is major, the observations are only separated by \(\sim 1\) d.

3.3 The soft branch

From observation #34, the rms is below \(\sim 5\) per cent and the count rate fades from \(\sim 1300\) to \(\sim 100\) count s\(^{-1}\). A linear relation can be observed, but much more scattered than during the HL (see Fig. 1). This new relation is followed by the system during the canonical soft state (M09).

3.4 Soft-to-hard transition: the adjacent hard line

Fig. 3 shows the RID in the log–log scale, where a monotonic rms increase is observed at constant flux (\(\sim 100\) count s\(^{-1}\)) from the bottom of the soft branch. Once the rms crosses the \(\sim 20\) per cent level, a linear decrease in rms and flux towards the HL is observed. During this adjacent HL (see note in Fig. 3), PDS are the typical properties of GX 339 during the 2002 outburst. They found that the optical and near-infrared properties of GX 339 are shown as the open triangles and open dots, respectively.

A sudden increase in rms is also observed during four observations. Three of these cross the \(\sim 10\) per cent line and show type-C QPOs, whereas a type-B QPO is observed for the other. Observations with type-A QPOs are not well differentiated from the soft branch. However, some of them are placed close or slightly above the 5 per cent rms line, which roughly delimits the soft branch (see also Fig. 2).

4 THE 2002 AND 2004 OUTBURSTS

In Figs 1–3, the RIDs corresponding to the 2002 and 2004 outbursts are shown as the open triangles and open dots, respectively.

(i) The 2002 RID is similar to the 2007 one (as in the case of the HID). The system leaves the HL roughly at the same position and an increase in flux at approximately constant (absolute) rms is also seen in the region where the accretion disc is observed by M09 during 2007. Homan et al. (2005) performed a multiwavelength campaign during the 2002 outburst. They found that the optical and near-infrared properties of GX 339–4 change on MJD 523 98, which corresponds to the first observation outside the HL (see Fig. 1). B05 also identified this observation as the first one belonging to the HIMS (i.e. non-LHS).

(ii) The 2004 outburst is fainter and therefore its absolute rms is also significantly lower. In this case, the variability does not decrease when the system leaves the HL, but increases with flux following a different linear relation. A clear feature related with the appearance of a thermal disc in the spectrum is not observed, suggesting a weaker disc component. Flux finally increases as rms decreases until the systems reach the soft branch. After a long permanence in this branch, rms increases up to \(\sim 20\) per cent and goes back to
the HL following the same adjacent HL observed in the other two outbursts.

The HL is identical for the three outbursts, independent of the maximum flux reached. We also observe this behaviour in the BHT H1743−322 (Muñoz-Darias et al., in preparation). The soft branch is also very similar between the three outbursts, lying within the ~1–5 per cent fractional rms region. We notice that observations just outside this region contain fast state transitions (see Nespoli et al. 2003; Casella et al. 2004). A detailed state classification in this region of the diagram is beyond the scope of this work.

It is also possible to define a region of the diagram (count rate $\geq 400 \text{ count s}^{-1}$ and $\sim 7$–10 per cent fractional rms), where only observations showing type-B QPOs (grey regions in Figs 1 and 2) are seen. These observations must be classified as SIMS according to Homan & Belloni (2005). We also identify two observations (70110-01-47.00 and 70108-03-02-00) with fainter type-B QPOs at a lower fractional rms than that of the region. A more detailed analysis of these data shows that the type-B QPO is only present in a part of the observation, with rms being that expected for the type-B QPO region.

5 HARD/SOFT RMS–INTENSITY DIAGRAMS

With the aim of tracing the evolution of soft and hard variability, we have computed the RID using soft (PCA 0–13; 2–6 keV) and hard (PCA 14–35; 6–15 keV) bands for the three outbursts. These energy ranges can be selected independently of the observing mode in which data were taken, which will allow us to cover the whole evolution of the three outbursts. In Fig. 4, the soft and hard RIDs are plotted for the 2007 outburst as the open circles and open triangles, respectively. It is clear that both RIDs show the same HL, where fractional rms drops from 40 to 30 per cent as the system becomes brighter. This is expected since it is well-known that rms is weakly energy-dependent (2–20 keV) during the LHS (see e.g. Muñoz-Darias et al., in preparation) for the cases of XTE J1752−388 (Muñoz-Darias et al. 2010) and Cyg X−1 (M09; see arrow in Fig. 4). For the same observations, an increase in flux at a constant rms is observed in the soft RID.

(ii) Since the disc component is detected by M09 (see arrows in Fig. 4), more fractional rms is seen systematically at higher energies, even at the very bottom of the soft branch.

On the other hand, the same adjacent HL is observed for both RIDs, as expected for LHS observations with flat rms spectra.

5.1 The 2002 and 2004 outbursts

The behaviour observed during the 2002 outburst (not shown here) is similar to the one described above. During the 2004 outburst, the vertical rise after leaving the HL is much less evident in the soft RID, suggesting a weaker accretion disc component. Neither the decrease in rms at a constant flux is observed in the hard RID, although the decay following the 20 per cent rms line is seen.

6 DISCUSSION

During the outburst evolution of BHT different accretions regimes are present. They yield the different states that we see when looking at the X-ray emission arising from these systems. In each one of these states, one or more physical mechanisms (e.g. thermal or non-thermal Comptonization in the corona, synchrotron emission from a jet, thermal emission from an accretion disc) play a role. The relative contribution of these mechanisms can be studied not only by using energy spectra, but also by analysing the particular variability they prompt in the light curves. In this work, we have shown that the evolution of the fast variability (0.1–64 Hz) can be used as a good tracer of the different accretion regimes in black hole binaries. This would have been impossible in case the variabilities associated with different spectral components were similar. Indeed, we find evidence for the presence of highly variable and an almost non-variable components in the light curves of GX 339−4:

(i) During the next five observations, the hard RID shows a decrease in rms at a constant count rate. Then, once the disc component appears in the spectrum (M09; see arrow in Fig. 4), a constant decrease in flux and rms following the 20 per cent line is seen until the type-B QPO appears (#33; star in Fig. 4). For the same observations, an increase in flux at a constant rms is observed in the soft RID.

(ii) Since the disc component is detected by M09 (see arrows in Fig. 4), more fractional rms is seen systematically at higher energies, even at the very bottom of the soft branch.

On the other hand, the same adjacent HL is observed for both RIDs, as expected for LHS observations with flat rms spectra.

(i) A sharp, linear relation between the absolute rms and flux is observed at the beginning and at the end the three outbursts of GX 339−4 we have studied. This has been observed before by Bleijsser et al. (2004) in Cyg X−1 during LHS observations. During these epochs, flux variations are probably tracing accretion rate variations and therefore rms is probably correlated with the accretion rate. The rms–flux relation was first observed by Uttley & McHardy (2001) at much shorter time-scales (~10 s) than the ones presented here (~1 ks). As a test we have studied the short-term variations in the rms–flux relation by dividing the light curve in 32-s segments for the first and last HL observations (2007 outburst). In both cases, we also found a linear relation, with the fractional rms being consistent with the average of the entire observation. We note that a detailed study on these short-term variations is beyond the scope of this work.

Uttley & McHardy (2001) found that the extrapolation of the rms–flux relation was consistent with no variability at lower count rates than the ones they studied. This supports the presence of two different components in the light curves. GX 339−4 offers the opportunity to study this linear relation over a much broader flux range. Consistent with the result found by Bleijsser et al. (2004) by extrapolating the relation they obtained for Cyg X−1, we do not see remaining flux when the rms is close to zero (see Fig. 3);
thus, our results are consistent with the presence of one component solely responsible for the variability at low count rates. Indeed, the system reaches its maximum variability level when this component dominates the spectrum. However, in the upper part of the HL, where Uttley & McHardy (2001) performed their study, the fractional rms decreases from 40 to 30 per cent. This could be associated with the presence of a second, less-variable component.

(ii) We find strong evidence for the presence of an extra component associated with the appearance of thermal emission in the energy spectrum of the source. This accretion disc component is not variable or varies very little. Its signature (i.e. increase in flux at constant rms) is clearly seen in the 2002 and 2007 outbursts, once the systems leave the HL. As expected, this becomes more evident if one only considers the 2–6 keV band (see Fig. 4). Indeed, The fractional variability of the system is lower at lower energies from the first observation in which the thermal component is detected, until it returns to the LHS at the end of the outburst.

(iii) The behaviour of the hard (6–14 keV) RID is complex. Once the system leaves the HL, a decrease in rms at constant flux is observed, suggesting that the properties of the physical mechanism that drives the hard emission are different. This could be explained by invoking a change in the properties of the corona or the base of the jet. For instance, the high-energy cut-off is observed to increase during the hard-to-soft transition (M09), which could be interpreted in terms of a cooling of the corona (see M09 and reference therein for details). It is remarkable that very few observations show fractional rms \( \lesssim 5 \) per cent in the hard band, suggesting that an underlying variable component (probably associated to power-law emission) is present even during the soft state (see Fig. 4).

The different rms–flux relations outlined above enable to perform a state classification of different observations by using the rms–intensity diagram. The following states can be identified using the description by Homan & Belloni (2005):

(i) HL observations correspond to the LHS, with rms in the range \( \sim 30–40 \) per cent. During these observations, PDS are qualitatively similar to each other, and can be fitted by using several broad Lorentzians (see e.g. fig. 3 in B05, where the three upper PDS correspond to observations at the beginning, the middle and the end of the HL during the 2002 outburst). We associated the sharp changes we see in the rms–flux relation at the end of the HL to the transition to the HIMS. This is supported by spectral and timing changes (e.g. B05; M09) and particularly by the coincidence in time between the drop we see in rms and that observed in the optical and infrared by Homan et al. (2003).

(ii) During the HIMS, the rms is observed to drop from \( \sim 30 \) to 10 per cent. The presence of the accretion disc in the spectrum is detected when the fractional rms is \( \sim 20 \) per cent. We note that this point marks the transition between the hard and the steep power-law or intermediate states according to McClintock & Remillard (2006) (also see M09). However, the major transition seems to occur before, when the system leaves the HL.

(iii) Soft-state observations are observed to follow a different rms–flux relation, consistent with the results obtained by Geisler et al. (2004) using relatively soft observations of Cyg X–1. The soft state in GX 339–4 is much softer than in Cyg X–1 and a soft branch very different from the HL is clearly seen in the RID. It shows a variability level lower than \( \sim 5 \) per cent and much more scatter than the HL. This suggests fast changes in (at least) one of the spectral components of the light curve.

(iv) We find that between \( \sim 7–10 \) per cent fractional rms, only type-B QPOs are observed (grey regions in Figs 1 and 2). They are SIMS observations according to Homan & Belloni (2005) and our diagram suggests that this kind of oscillations could follow a particular rms–flux relation. We note that transitions between soft and SIMS are very fast, which result in hybrid observations where timing properties are not constant. Thus, an accurate state classification close to the \( \sim 5 \) per cent rms line is in some cases not compatible with using entire 1–3 ks observations.

(v) An adjacent HL is observed at the end of the three outbursts. Observations during this adjacent HL show PDS typical of the LHS but with a slightly softer spectrum than that of the HL. It is unclear why this other linear relation, which is not seen when the system leaves the HL, is present. It could suggest the presence of new components (see Russell et al. 2010) or the total absence of any of the components (e.g. accretion disc) observed during the flux-rise phase at the beginning of the outburst.

7 CONCLUSIONS

We have proved that the evolution of the fast variability can be used as a good tracer of different accretion regimes in black hole binaries. In particular, we find that apart from the linear rms–flux relation found during LHS, different relations are followed by black hole binaries during the soft and intermediate states.

In this work, we have presented the rms–intensity diagram, showing that it is possible to associate different regions of this diagram to the different states observed in black hole transients. Transitions can also be identified, thanks to the marked changes they produce in the diagram. The application of this new tool to other systems and its interpretation will provide new insights into our understanding of the physical processes that take place during accretion episodes in X-ray binaries.

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