Probabilistic naturalness measure for dipole moments due to new physics

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We introduce a probabilistic measure of naturalness (naturalness level) to fix naturalness bounds quantitatively. It is applied to the anomalous magnetic moments and the electric dipole moments due to new physics.

PACS number(s): 12.60.-i, 02.50.Cw, 11.10.Gh, 13.40.Em

Quest for “naturalness” [1][2], without fine-tuning cancellations motivates and promotes the people to develop beyond-the-standard model physics, such as composite [3], grand unification [4], supersymmetry [5], superstring [6], braneworlds [7], and so on. The naturalness requirements, in addition to the experimental tests, serves as a guiding principle for new physics model building, and bring us useful information on their parameters [8][9][10]. In the previous paper [10], we applied it to the naturalness of quark-lepton mass renormalization, and derived many model-independent bounds on their anomalous magnetic moments and electric dipole moments. The naturalness relations, however, have so far been ambiguous order-of-magnitude relations, and have had no common basis to compare them model-to-model. In this paper, we would like to propose a universal probabilistic measure of naturalness, which we call “naturalness level”. We apply them to the previously derived naturalness bounds [10] on the dipole moments. For a given definite naturalness level, we can assign more definite values for the bounds than the previous ones.

Let us consider what we usually mean by naturalness in physics. For example, if the renormalized mass

\[ m' = m^b + \delta m \]  (1)

is very small compared with the bare mass \( m^b \) and its quantum correction \( \delta m \), we say that it is unnatural fine-tuning cancellation. In saying so, we are expecting independent and random appearance of \( m^b \) and \( \delta m \), neglecting the further information at this stage. If we know the true mechanism, it would no longer be unnatural. In general, when we say something natural or unnatural, we seem always to assume some randomness at some non-ultimate stage, and if the probability based on the randomness is small we say it unnatural. Though the cancellations could have complicated form in general, the quantity should be defined by homogeneous function of its contributions for meaningful comparison of their magnitudes. Thus we formulate it as follows.

Suppose, at some stage in some model, a quantity \( X \) is given by

\[ X = f(X_1, X_2, \ldots) \]  (2)

where \( f \) is a homogeneous function (i.e. \( \forall a, aX = f(aX_1, aX_2, \ldots) \)) of the contributions \( X_1, X_2, \ldots \). For example, for (1), we take \( X = m', X_1 = m^b \) and \( X_2 = \delta m \). Then we assume that the \( X_i \)'s appear with uniform randomness in \( X_i \)-space. The naturalness means that \( |X| \) should not be too small in comparison with the individual contributions \( X_i \), or, in general, with some homogeneous combination of them \( g(X_1, X_2, \ldots) (\forall a, ag(X_1, X_2, \ldots) = g(aX_1, aX_2, \ldots)) \). Then, the naturalness bound is expressed as

\[ |X| > \epsilon|g(X_1, X_2, \ldots)| \quad \text{(naturalness bound)} \]  (3)

with some given small positive constant \( \epsilon \). For example, for (1), \( |m'| > \epsilon|\delta m| \), \( |m'| > \epsilon \sqrt{(m^b)^2 + (\delta m)^2} \), etc.. Then we define the “naturalness level (NL)” \( p \) of the naturalness bound (3) as its “probability” based on the assumed randomness:

\[ p = P(|X| > \epsilon|g(X_1, X_2, \ldots)|), \quad \text{(naturalness level)} \]  (4)

which would be the statistical confidence level of the inequality (3), if the “probability” be real probability. Because we assume uniform randomness of \( X_i \), the naturalness level \( p \) is the relative volume occupied by the region of (3) in the \( X_i \)-space. Since the functions \( f \) and \( g \) are homogeneous, it is sufficient to evaluate the relative volume of the region (3) on the hypersurface \( \sum |X_i|^2 = 1 \).

In the following, we derive the expressions for the naturalness level \( p \) in terms of \( \epsilon \) for the naturalness bounds in lepton- and quark-mass renormalizations. We invert the relation to express \( \epsilon \) in terms of \( p \), which enables us to fix the values of the naturalness bounds for given naturalness levels. Then, we apply them to examples of the naturalness bounds on the anomalous magnetic moments and the electric dipole moments of leptons and quarks due to new physics [11].

(A) Mass Renormalization under conserved CP

In this case, the renormalized mass takes the form in (4) above. Let us consider the naturalness bound

\[ |m'| > \epsilon|\delta m|. \]  (5)

In terms of the polar angle \( \phi = \tan^{-1}(X_2/X_1) = \tan^{-1}(\delta m/m^b) \) in the \( X_1 \)-\( X_2 \) plane, (3) is rewritten as
By the general arguments above, the naturalness level \( p \) is \( l/2\pi \) with the total length \( l \) of the arc allowed by (10) on the unit circle \( X_1^2 + X_2^2 = 1 \), and is given by

\[
p = \frac{1}{\pi} \cot^{-1} \frac{\epsilon^2 - 2}{2\epsilon}.
\]

(7)

Note that for \( \epsilon \to 0(\infty), p \to 1(0) \). For small \( \epsilon \),

\[
p = 1 - \frac{\epsilon}{\pi} - \frac{\epsilon^3}{6\pi} + O(\epsilon^5).
\]

(8)

Eq. (7) is inverted for \( \epsilon \) as (with \( q = 1 - p \))

\[
\epsilon = \sqrt{2 + \cot^2 p \pi + \cot p \pi \tau = \pi q - \frac{\pi^3 \tau^3}{6} + O(q^5)}.
\]

(9)

This means that at NL (naturalness level) = 90%(95%),

\[
|m| > 0.309|\delta m| \quad (|m| > 0.156|\delta m|).
\]

(10)

Now we apply this to the effects of the anomalous magnetic moment \( \mu \) due to new physics \( [1] \). The contribution is given by

\[
\delta m = -3eQ\mu \Lambda^2/8\pi^2,
\]

(11)

where \( e \) is the electric coupling constant, \( Q \) is the electric charge of the fermion, and \( \Lambda \) is the new physics scale. In practice, \( \Lambda \) is the photon momentum cutoff of the self mass diagram with the cutoff factor \((1 - k^2/\Lambda^2)^{-2} \) where \( k \) is the momentum of the photon propagator. The lepton scattering experiments sets the bound

\[
\Lambda = (2\pi)^{-1/2} e\Lambda^{VV} > 2.26\text{ TeV},
\]

(12)

where we have used the phenomenological bound on the contact-interaction scale \( \Lambda^{VV} = 18.0\text{ TeV} \) with vector-vector coupling and with negative sign \( [12] \). (We use the value \( \epsilon = 0.315 \) at the mass scale of 200 GeV.)

Using \( [12] \) and the experimental values for the masses \( [13] \), we obtain the bounds for \( \delta a = \mu/(e_0\Lambda^{2m}) \) \((e_0/4\pi = 1/137)\), at NL=90%(95%)

\[
|\delta a_e| < 9.1 \times 10^{-11} (1.8 \times 10^{-10}), \\
|\delta a_{\mu}| < 3.9 \times 10^{-12} (7.7 \times 10^{-11}), \\
|\delta a_{\tau}| < 1.1 \times 10^{-3} (2.2 \times 10^{-3}).
\]

(13)

The subscripts of \( \delta a, \mu \) and \( d \) here and below indicate fermion species.

The recent experimental result on \( \delta a_{\mu} \) \( [14] \), combined with the standard-model prediction \( [13] \), set the 95% confidence level lower bound

\[
\delta a_{\mu} > 5.2 \times 10^{-10}.
\]

(14)

Combining \( [1], [11] \) and \( [14] \) we obtain the upper bound for the new physics scale \( \Lambda \) at NL=90%(95%)

\[
\Lambda < 200\text{ TeV} \quad (280\text{ TeV}).
\]

(15)

For quarks, we use phenomenological values of the masses \( [13] \) to obtain the bounds at NL=90%(95%)

\[
|\mu_1|/|\mu_N| < 2.2 \times 10^{-6} (4.4 \times 10^{-6}), \\
|\mu_2|/|\mu_N| < 8.4 \times 10^{-6} (1.7 \times 10^{-5}), \\
|\mu_3|/|\mu_N| < 1.5 \times 10^{-4} (3.0 \times 10^{-4}), \\
|\mu_4|/|\mu_N| < 6.9 \times 10^{-4} (1.4 \times 10^{-3}), \\
|\mu_5|/|\mu_N| < 4.4 \times 10^{-3} (8.8 \times 10^{-3}),
\]

(16)

where \( \mu_N \) is the nuclear magneton.

For neutrinos, instead of (11) we have

\[
\delta m = -3e^2g^2c\mu \Lambda^2/64\pi^4,
\]

(17)

where \( g \) is the gauge coupling constant of the electroweak gauge group SU(2)_L, and \( c \) is a numerical constant of \( O(1) \), which is in principle calculable. It requires, however, lengthy calculations and careful treatments of divergent integrals, and is presently under investigation. Apart from \( c \), we obtain, at NL=90%(95%),

\[
|\mu_{\nu_e}|/|\mu_B| < 9.9 \times 10^{-14} (2.0 \times 10^{-13}), \\
|\mu_{\nu_\mu}|/|\mu_B| < 6.2 \times 10^{-9} (1.2 \times 10^{-8}),
\]

(18)

where \( \mu_B \) is the Bohr magneton.

(B) **Mass Renormalization with broken CP**

When CP invariance is broken, the fermion mass have the 1- and \( \gamma_5 \)- components (indicated by the subscripts 0 and 5, respectively),

\[
m^0 = m^0_b + \delta m_0, \quad m^5 = m^5_b + \delta m_5.
\]

(19)

The physical mass is given by the expression

\[
m^r = \sqrt{(m^0_b + \delta m_0)^2 + (m^5_b + \delta m_5)^2}.
\]

(20)

which is taken as the \( X \) in the above general arguments. It is homogeneous in \( X_1 = m^0_b, X_2 = \delta m_0, X_3 = m^5_b, \) and \( X_4 = \delta m_5 \), which are taken to be uniformly random. Let us consider the naturalness bound

\[
|m^r| > \epsilon|\delta m_0|, \quad |m^r| > \epsilon|\delta m_5|.
\]

(21)

The naturalness level \( p \) is given by \( V/2\pi^2 \), where \( V \) is the volume of the part allowed by (21) on the hypersphere \( \sum_{i=1}^4 |X_i|^2 = 1 \). A lengthy but straightforward calculation leads to the expression

\[
p = 1 - \frac{2}{\pi} \int_0^\phi \sqrt{\alpha - \sin^2 \phi} d\phi,
\]

(22)

where \( \sin^2 \phi = \alpha \) and

\[
\alpha = \left[ 1 - (1 - \epsilon^2/2)\sqrt{1 + \epsilon^4/4} \right]/2,
\]
\[ \beta = \left[ 1 + \frac{\epsilon^2}{2} \right] / \sqrt{1 + \frac{\epsilon^2}{4}}. \]  
(23)

For small \( \epsilon \)

\[ p = 1 - \frac{\epsilon^2}{8} \left( 1 + \frac{3\epsilon^2}{16} - \frac{5\epsilon^4}{64} + \cdots \right). \]  
(24)

It is inverted as (with \( q = 1 - p \))

\[ \epsilon = 2\sqrt{2q} \left( 1 - \frac{3q}{4} + \frac{143q^2}{32} + \cdots \right). \]  
(25)

This means that at NL 90%(95%),

\[ |m'\rangle > 0.867|\delta m_0\rangle \quad \text{or} \quad |m'\rangle > 0.616|\delta m_0\rangle, \]
\[ |m'\rangle > 0.867|\delta m_5\rangle \quad \text{or} \quad |m'\rangle > 0.616|\delta m_5\rangle. \]  
(26)

Note that, for small \( p\), \( p \sim O(\epsilon) \) in (A), while \( p \sim O(\epsilon^2) \) in (B), so that the corresponding bounds in (24) are more stringent than those in (11).

Now we apply them to the effects of the anomalous magnetic moment \( \mu \) and the electric dipole moment \( d \) due to new physics [1]. The contribution \( \delta m_0 \) is the same as the \( \delta m \) in (11), while \( \delta m_5 \) is given by

\[ \delta m_5 = -3eQd\Lambda^2/8\pi. \]  
(27)

Then, we obtain at NL=90%(95%)

\[ |\delta \alpha_e| < 3.3 \times 10^{-11} \quad (4.6 \times 10^{-11}), \]
\[ |\delta \mu_\mu| < 1.4 \times 10^{-6} \quad (2.0 \times 10^{-6}), \]
\[ |\delta \alpha_\tau| < 3.9 \times 10^{-4} \quad (5.5 \times 10^{-4}), \]
\[ |d_e| < 6.3 \times 10^{-22} \quad (8.8 \times 10^{-22}) \text{ecm}, \]
\[ |d_\mu| < 1.3 \times 10^{-19} \quad (1.8 \times 10^{-19}) \text{ecm}, \]
\[ |d_\tau| < 2.2 \times 10^{-18} \quad (3.1 \times 10^{-18}) \text{ecm}. \]  
(28)

The upper bound for the new physics scale \( \Lambda \) from [21], [1] and [1] at NL=90%(95%) is

\[ \Lambda < 120\text{TeV} \quad (140\text{TeV}). \]  
(29)

Note that bounds on \( \delta \alpha 's \) in (28) and that on \( \Lambda \) in (29) are more stringent than those in [13] and [14]. In general an additional condition (here the CP conservation) makes the bounds less stringent.

For quarks, we obtain the bounds at NL=90%(95%)

\[ |\mu_\alpha|/\mu_N < 7.9 \times 10^{-7} \quad (1.1 \times 10^{-6}), \]
\[ |\mu_d|/\mu_N < 3.0 \times 10^{-6} \quad (4.2 \times 10^{-6}), \]
\[ |\mu_s|/\mu_N < 5.4 \times 10^{-5} \quad (7.7 \times 10^{-5}), \]
\[ |\mu_c|/\mu_N < 2.5 \times 10^{-4} \quad (3.5 \times 10^{-4}), \]
\[ |\mu_t|/\mu_N < 1.6 \times 10^{-3} \quad (2.2 \times 10^{-3}), \]
\[ |\mu_\tau|/\mu_N < 3.2 \times 10^{-2} \quad (4.6 \times 10^{-2}), \]
\[ |d_\alpha| < 8.3 \times 10^{-21} \quad (1.2 \times 10^{-20}) \text{ecm}, \]
\[ |d_d| < 3.1 \times 10^{-20} \quad (4.4 \times 10^{-20}) \text{ecm}, \]
\[ |d_s| < 5.7 \times 10^{-19} \quad (8.1 \times 10^{-19}) \text{ecm}, \]
\[ |d_c| < 2.6 \times 10^{-18} \quad (3.6 \times 10^{-18}) \text{ecm}, \]
\[ |d_t| < 1.7 \times 10^{-17} \quad (2.3 \times 10^{-17}) \text{ecm}, \]
\[ |d_\tau| < 3.4 \times 10^{-16} \quad (4.8 \times 10^{-16}) \text{ecm}. \]  
(30)

For neutrinos, \( \delta m_0 \) is the same as the \( \delta m \) in (17), while \( \delta m_5 \) is given by

\[ \delta m_5 = -3eQd\Lambda^2/64\pi^4. \]  
(31)

Then, we obtain, at NL=90%(95%),

\[ |\mu_\alpha|/\mu_B < 3.5 \times 10^{-14} \quad (5.0 \times 10^{-14}), \]
\[ |\mu_d|/\mu_B < 2.2 \times 10^{-9} \quad (3.1 \times 10^{-9}), \]
\[ |\mu_s|/\mu_B < 2.1 \times 10^{-7} \quad (3.0 \times 10^{-7}), \]
\[ |d_\alpha| < 6.8 \times 10^{-20} \quad (9.6 \times 10^{-20}) \text{ecm}, \]
\[ |d_d| < 4.3 \times 10^{-20} \quad (6.1 \times 10^{-20}) \text{ecm}, \]
\[ |d_s| < 4.1 \times 10^{-18} \quad (5.8 \times 10^{-18}) \text{ecm}. \]  
(32)

The presently available experimental bounds are

\[ \delta \alpha_e = (1.2 \pm 2.8) \times 10^{-11}, \]
\[ d_e = (6.9 \pm 7.4) \times 10^{-28} \text{ecm}, \]
\[ d_\mu = (26 \pm 10) \times 10^{-11}, \]
\[ d_\tau = (3.7 \pm 3.4) \times 10^{-10} \text{ecm}, \]
\[ -0.052 < \delta \alpha_e < 0.058 \quad (95\% \text{CL}), \]
\[ -2.2 < R e d_\tau < 4.5 \times 10^{-17} \text{ecm} \quad (95\% \text{CL}), \]
\[ -2.5 < I m d_\tau < 0.8 \times 10^{-17} \text{ecm} \quad (95\% \text{CL}), \]
\[ |\mu_\alpha| < 1.5 \times 10^{-10} \text{H}_B \quad (90\% \text{CL}), \]
\[ |\mu_\mu| < 6.8 \times 10^{-10} \text{H}_B \quad (90\% \text{CL}), \]
\[ |\mu_\tau| < 3.9 \times 10^{-7} \text{H}_B \quad (90\% \text{CL}), \]
\[ |d_\alpha| < 5.2 \times 10^{-17} \text{ecm} \quad (95\% \text{CL}). \]  
(33)

The naturalness bounds for \( \delta \alpha, d, \mu, \tau, \mu_e, \text{and} \ d_\nu \) are more stringent than the experimental bounds, while those for \( \delta \alpha_e, \delta \mu_\mu, \delta_\tau \) and \( \mu_\nu \) are less stringent, and that for \( \mu_\nu \) is comparable. If we use the information on the neutrino mass differences from the experimental results on the solar [21] and atmospheric [25] neutrinos, the naturalness bounds for all the neutrinos become the same as those for \( \nu_e \) in (18) and (33), which are much more stringent than experimental ones.

We can apply the naturalness-level analysis to various problems. For example, we can calculate the naturalness levels for the mass ratios \( m_\alpha/m_\mu, m_\mu/m_\tau, m_\alpha/m_\tau \), etc. and the mixing angles and the CP violating phase, under the assumption that they are derived by the diagonalization of random mass matrices. We can calculate the naturalness levels and compare them under the various naturalness-achieving mechanisms, such as seesaw mechanism [21], compositeness [1], supersymmetry [2], braneworld etc [5]. Some of them are presently under investigations, and will be reported elsewhere.

The naturalness level defined here possesses very common quantitative meaning independent of specific quantities in specific theories. It depends, however, on the choice of the “stage” where we assume the randomness of the contributions to the quantity. The lower naturalness requires the further theoretical explanations from the more fundamental stage so as to improve the naturalness level. We expect that such naturalness-level analyses
would provide powerful quantitative guides in search for true explanations of the nature.

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