Blunt trauma to large vessels: a mathematical study

Rovshan M Ismailov*1, Nikolai A Shevchuk2,3, Joseph Schwerha4, Lawrence Keller4 and Higmat Khusanov5

Address: 1Department of Epidemiology, Graduate School of Public Health, University of Pittsburgh, Pittsburgh, PA 15213, USA, 2Center for Cancer and Immunology Research, Children’s Research Institute, Washington, DC, USA, 3Institute for Biomedical Sciences/Program in Molecular and Cellular Oncology, Washington, DC, USA, 4Department of Environmental and Occupational Health, Graduate School of Public Health, University of Pittsburgh, Pittsburgh, PA 15213, USA and 5Institute of Mechanics and Seismic Stability of Structures, Academy of Science of Uzbekistan, Tashkent, Uzbekistan

Email: Rovshan M Ismailov* - rmi1@pitt.edu; Nikolai A Shevchuk - niash@gwu.edu; Joseph Schwerha - jshwerha@ceoh.pitt.edu; Lawrence Keller - lkeller@ceoh.pitt.edu; Higmat Khusanov - instmech@uzsci.net

* Corresponding author

Abstract

Background: Blunt trauma causes short-term compression of some or all parts of the chest, abdomen or pelvis and changes hemodynamics of the blood. Short-term compression caused by trauma also results in a short-term decrease in the diameter of blood vessels. It has been shown that with a sudden change in the diameter of a tube or in the direction of the flow, the slower-moving fluid near the wall stops or reverses direction, which is known as boundary layer separation (BLS). We hypothesized that a sudden change in the diameter of elastic vessel that results from compression may lead not only to BLS but also to other hemodynamic changes that can damage endothelium.

Methods: We applied Navier-Stokes, multiphase and boundary layer equations to examine such stress. The method of approximation to solve the BL equations was used. Experiments were conducted in an aerodynamic tube, where incident flow velocity and weight of carriage with particles before and after blowing were measured.

Results: We found that sudden compression resulting from trauma leads to (1) BLS on the curved surface of the vessel wall; (2) transfer of laminar boundary layer into turbulent boundary layer. Damage to the endothelium can occur if compression is at least 25% and velocity is greater than 2.4 m/s or if compression is at least 10% and velocity is greater than 2.9 m/s.

Conclusion: Our research may point up new ways of reducing the damage from blunt trauma to large vessels. It has the potential for improvement of safety features of motor vehicles. This work will better our understanding of the precise mechanics and critical variables involved in diagnosis and prevention of blunt trauma to large vessels.

Background

Blunt trauma to large vessels is well recognized outcome in patients with chest, abdominal or pelvic contusion [1,2]. Thoracic aortic rupture carries a mortality rate of over 90%, with on-scene death occurring in more than 80% of individuals with this injury [1,2]. Even minimal aortic injuries are important, since the potential adverse consequences of not operatively repairing minimal aortic...
injuries, defined as a small (<1 cm) intimae flap with no or minimal periaortic hematoma, include formation, enlargement, and rupture of pseudoaneurysm, embolism of loose intima, or thrombus and progressive dissection of the aortic wall [3,4].

Compression of the chest, abdomen or pelvis as a result of injury changes hemodynamics of the blood. This short-term compression also causes a transient decrease in the diameter of the aorta and other blood vessels. Fox and McDonald have shown that a sudden change in the diameter of the tube or in the direction of the flow, causes the slower-moving fluid near the wall to lose momentum and stop or reverse direction [5]. In such situations, the core flow is then separated from the wall by a fluid disturbance or separation zone [5]. Boundary layer separation (BLS) has also been shown to appear due to adverse pressure gradients (i.e. pressure increasing in the direction of the flow) [6]. It is important to note, however, that BLS does not take place in an elastic vessel under normal physiological conditions [7].

In this work, we tested the hypothesis that a sudden change in the diameter of an elastic vessel that results from compression can cause not only BLS but also other hemodynamic changes that cause large vessel injury via damage to the endothelium. We further explored and analyzed shear stress at the compressed part of the vessel and conducted a series of experiments to explore certain hemodynamic changes upon compression caused by injury.

Materials and methods
All experiments were conducted in aerodynamic tube (AT) with a square section (200 × 200 mm²). The experimental part was located at a distance of 6.0 m from the rectifying lattices. In order to shape the flat laminar boundary layer (LBL), a polished steel plate with sharp edges was used. The upper part of the plate was combined and had carriage with the particles. To avoid BLS, the transverse joints of the walls of a carriage and separate steel plate were thoroughly condensed and painted over. A steel plate was placed into the wind tunnel and was blown during the specific time. The specific goals of the experiments were: (1) to measure incident flow velocity and (2) to determine the weight of the carriage with particles before and after blowing. In all experiments, the degree of turbulency of incident flow was approximately the same and equal to 0.13%.

Compression caused by trauma leads to stress (stress is defined here as the force applied to a certain surface). We applied fundamental differential equations, namely Navier-Stokes, multiphase and boundary layer (BL) equations, to examine such stress and evaluate the related hemodynamics as well.

In the equations below (unless otherwise specified) we will assume that $\rho$ is blood density; $u$ is longitudinal blood stream velocity near the vessel wall (inside the BL); $v$ is transverse blood stream velocity near the vessel wall (inside the BL); $U$ is longitudinal blood stream velocity in the center of the vessel; $y$ is transverse blood stream velocity in the center of the vessel (see figure 1); $p$ – pressure; $\mu$ – viscosity; $t$ – time; $\tau$ – shear stress; $\omega$ – acceleration.

**BLS resulting from compression caused by injury**

It is well known that a sudden change in vessel diameter leads to BLS [8]. This follows from an equation for BL and boundary conditions:

If $u = 0, v = 0$ and $y = 0$, then

$$\mu \left( \frac{d^2 u}{d^2 y} \right)_{y=0} = \frac{dp}{dx}.$$  \hspace{1cm} (2.1.1)

If pressure is decreasing (accelerated flow), then $\frac{dp}{dx} < 0$ and from (2.1.1) it follows that

$$\frac{d^2 u}{dy^2} < 0.$$

If pressure is increasing (decelerating flow) then: $\frac{dp}{dx} > 0$ and from (2.1.1) it follows that

$$\frac{d^2 u}{dy^2} > 0.$$

There is also inside the BL

$$\frac{d^2 u}{dy^2} = 0.$$  \hspace{1cm} (2.1.2)

Equation (2.1.2) describes the boundary between the direct and recurrent flow, which also satisfies the following condition $\left( \frac{du}{dy} \right)_{y=0} = 0$.

Physically it can be described as follows: A flow that is "broken" in the BL has so little kinetic energy that it cannot move further toward the high-pressure area so this flow moves away from the vessel wall and toward the axial (central) flow. Thus, injury creates conditions for BLS. Note that there is no BLS under normal physiological conditions.
conditions, although certain events can cause such separation [9].

Compression can create conditions for accelerated blood flow movement along a relatively curved surface of the vessel and along a relatively flat surface of the vessel. Let’s examine the motion on the curved surface. The similarity law is used. A mobile BL is formed when the body moves with acceleration and vessel surface is covered by accelerated flow. Schlichting has provided calculations for accelerated motion of a body in the standing fluid [8]. According to these calculations, within the range of Reynolds number of 2,000 – 20,000, time necessary for BLS changes from 0.039 to 0.392 s.

Shear stress at sudden and accelerated blood flow
Theoretical studies [8] have shown that the length of time necessary for BLS depends on the type of acceleration. According to Schlihting [8], motion can be sudden (type I) and accelerated (type II).

Taking into account the relativity of motion (identical \( \frac{du}{dx} \)), we have the following formula for type I (sudden acceleration): detachment

\[
1 + \left(1 + \frac{4}{3\pi} \left( \frac{du}{dx} \right) \right) t_{\text{attachment1}} = 0 \tag{2.2.1}
\]

For type II (moderate acceleration):

\[
1 + 0.427 \left( \frac{du}{dx} \right) t_{\text{attachment2}} = 0 \tag{2.2.2}
\]

where \( t_{\text{attachment}} \) denotes the point that removes the BL (it appears in special cases examined below).

Experiments inside the LBL
Let’s examine the motion on the flat surface in the LBL. Let’s establish the connection between the damage to the endothelium (removal of particles from the vessel wall [7] and shear stress. Taking into account the similarity law (mechanically similar flows) namely the ratio of inertial forces to the frictional forces, an AT can be used to study the damage inside the BL within the range of Reynolds number. For the Reynolds number of the canine ascending aorta [7], the peak value for the velocity is: 1,587 < \( \text{Re} \) < 11,500 and for human ascending aorta [10]: 2,645 < \( \text{Re} \) < 12,566.

In our experiments, the anterior section of the flat surface with the sharp leading edge corresponds to LBL. The overall critical length of LBL is determined by Reynolds number:

\[
\text{Re} = \left( \frac{U_{x}}{\nu} \right) < 5 \cdot 10^5 + 10^6 \tag{2.3.1}
\]
where x is a distance from the anterior edge of the plate; v is kinematic viscosity; $u_\infty$ is incident flow velocity.

Experiments in AT were conducted at the following range of Reynolds number:

$$1300 \leq Re \leq 100,000$$

which is considerably less than the critical number (2.3.1) for the flat surface.

**A mathematical model of the damage to the wall**

For the more detailed analysis, let’s examine a mathematical model. To describe the above process, two-phase model is used [11]. Generally, this model is used to describe blood flow in the large arteries or filtration of fluid through the capillary wall [12]. Polyphase models are based on the idea of interacting and interpenetrating media [11]. The system of particles (for the two-phase flow) in this model is replaced by mathematical continuum [11]. In addition, it is assumed that particle size is considerably less than the distance at which flow conditions may change [11].

**Damage at accelerated flow**

Let’s examine damage caused by accelerating motion. To examine the LBL at an accelerated flow, experiments were conducted using the methods discussed above except that the circumfluent steel plate was added to create a fast delayed flow.

**Experiments inside the TBL**

We conducted experiments in the developed TBL in the range of Reynolds number (for AT): $70,000 < Re < 140,000$

**Results**

**BLS resulting from compression caused by injury**

Within the range of Reynolds number of $2,000 – 20,000$, time necessary for BLS changes from 0.039 to 0.392 s [8]. It is well known that peak velocity of blood flow in canine aorta changes within the range of Reynolds number 1,587 to 11,500 and in human aorta from 2,645 to 12,566 [10].

Theoretically, injury causing sudden compression of the aorta can happen within the range of Reynolds number of $2,000 – 20,000$ and within 0.039 to 0.392 s. Therefore, injury that causes sudden compression may result in BLS. Note that there is no BLS under normal physiological conditions, although certain events can cause such separation [9].

**Shear stress at sudden and accelerated blood flow**

If we compare (2.2.1) and (2.2.2) and take into account that $du / dy$ is identical, then the time necessary for BLS is 3.3 times longer with the second type of motion as follows from the equation: $1.424 \ t_{\text{detachment}1} = 0.427 \ t_{\text{detachment}2}$

Let’s examine shear stress for these 2 types of motion.

The relation of type I to type II stress is

$$\frac{\tau_1}{\tau_2} = \sqrt{\left(\frac{\mu \rho (u_1)^3}{l_1}\right) / \sqrt{\left(\mu \rho (u_2)^3 / l_2\right)}} ,$$

and taking into account that the motion is accelerated

$$\frac{\tau_1}{\tau_2} = 3 \sqrt{\left(\frac{u_1}{u_2}\right)} ,$$

Therefore, shear stress appearing on the vessel wall for the first type of acceleration (sudden) is 3 times higher than the stress caused by moderate acceleration. Thus, the value of shear stress depends on the type of acceleration.

Let’s examine shear stress at both types of motion can exceed the critical value that can damage the endothelium. If Pausel type of fluid flow is taken into account, shear stress can be determined according to formula

$$\tau = \frac{8\mu u}{d} [7]$$

where $\mu$ – fluid viscosity, $u$ – average velocity, $d$ – vessel diameter.

It is known that the average normal shear stress is $0.43 \cdot \frac{N}{m^2}$ for the canine ascending aorta, $0.53 \cdot \frac{N}{m^2}$ for the canine abdominal aorta and $0.8 \cdot \frac{N}{m^2}$ for the canine femoral artery [7,8]. In addition, it has been found that shear stress exceeding $40 \cdot \frac{N}{m^2}$ can damage the layer of endothelial cells [13]. If canine aorta was compressed by 30% as the result of injury and acceleration was moderate, then, according to formula [7]:

$$d = 0.015m, \ \alpha = 13.2, \ \mu = 1.2 \frac{m}{s}, \ \varphi = 0.004 \frac{N \cdot s}{m^2}, \ \tau_1 = 13.88 \frac{N}{m^2}$$

Where $\alpha$ is Womersley parameter [7].

However, under the same conditions, but with sudden compression, according to formula (2.2.1):

$$\tau_2 \geq 41.64 \frac{N}{m^2}$$

As it was mentioned above, shear stress exceeding $40 \cdot \frac{N}{m^2}$ may not only cause damage to the endothelium but also lead to atheromatosis [14]. Acceleration in this case cannot initialize BLS. Therefore, sudden compression caused by injury is detrimental if it happened within the
range of Reynolds numbers specified above and within the corresponding time interval.

**Experiments inside the LBL**

Assuming that there is a link between shear stress and the damage to the endothelium, the stress must be decreasing downstream since shear stress is the decreasing function from x:

\[ \tau = \mu U_\infty \alpha \sqrt{U_\infty / \nu x}. \]

Where \( \alpha \) is a coefficient of local shear stress [8].

Figure 2 provides the summary of results. According to this figure, the damage is influenced by the incident flow velocity (the greater the velocity the greater shear stress (along the longitudinal axis)). Therefore, the damage is a function of shear stress.

**A mathematical model of the damage**

Let's assume that carrying flow, which is going parallel to the flat surface of the vessel, is viscous, and the flow itself (solid particles (erythrocytes)) is ideal. Also we will assume that axis x is along the flow, while axis y is perpendicular to the surface. Then two equations for plasma (1st phase) will take the following form:

\[ \begin{align*}
\rho u_1 \frac{\partial u_1}{\partial x} + \rho v_1 \frac{\partial u_1}{\partial y} &= \mu \frac{\partial^2 u_1}{\partial y^2} + k(u_2 - u_1), \\
\rho \frac{\partial u_2}{\partial x} + \rho v_2 \frac{\partial u_2}{\partial y} &= k(u_2 - u_1), \\
\rho \frac{\partial v_2}{\partial x} + \rho v_2 \frac{\partial v_2}{\partial y} &= k(v_2 - v_1), \\
\frac{\partial \rho_2 u_2}{\partial x} + \frac{\partial \rho_2 v_2}{\partial y} &= 0.
\end{align*} \]  

And an equation for 1st and 2nd phases will take the following form:

\[ \begin{align*}
\rho_2 + \rho_1 &= 1. \\
\frac{\partial u_1}{\partial x} + \frac{\partial v_1}{\partial y} &= \rho \frac{\partial u_1}{\partial x} + \rho \frac{\partial v_1}{\partial y} \\
\rho_1 \frac{\partial \rho_2}{\partial x} + \rho_2 \frac{\partial \rho_1}{\partial y} &= f_1(\tau_1), \\
\rho_1 &= 0; \\
\rho_2 &= \rho_2^*; \\
f_1 &= 0.
\end{align*} \]  

Written above \( u_1 \) and \( u_2 \) are longitudinal components of velocity; \( v_1 \) and \( v_2 \) are transverse components of velocity; \( f_1 \) and \( f_2 \) – initial distribution of the velocities in the BL; \( \rho_2^* \) and \( \rho_1^* \) initial distribution of the densities. The beginning of the 2nd phase (due to the removal) was determined according to the calculated shear stress and the corresponding removal according to the experimental data.

The system of equations (3.4.1) – (3.4.6) with boundary conditions (3.4.7) – (3.4.9) was solved numerically by means of application of the finite-difference method with order of approximation equal to 0.0001.

Figure 3 reflects the distribution of the transverse velocity of the first and second phase. According to this figure, the denser phase moves to the wall, carrying the flow to the
wall and causing an additional mixture of the main flow with the wall flow. This denser phase creates additional tension on the wall unless we assume that the second phase is an ideal flow.

Since we assume the carrying flow to be viscous and the flow with erythrocytes viscous as well, then viscosity itself will be determined according to formula [15]:

\[
\mu_2 = \mu_1 \frac{\theta_1}{\theta_2} - 0.39 \frac{dp}{dx} \frac{R^4}{\theta_2}
\]

where \(\theta_1\) is the expenditure of the carrying phase; \(\theta_2\) is the expenditure of denser phase; \(\mu_1\) is a coefficient of dynamic viscosity; \(\frac{dp}{dx}\) is pressure gradient, \(R\) is a radius of a pipe. Shear stress will be then determined by the formula: \(\tau_2 = \mu_2 f_0 \frac{\partial u_2}{\partial y}\), where \(f_0\) is a dimensionless velocity.

The system of equations (3.4.1) – (3.4.6) will be changed: instead of the equations (3.4.3) and (3.4.4) there are will be:

\[
\rho_2 u_2 \frac{\partial u_2}{\partial x} + \rho_2 v_2 \frac{\partial u_2}{\partial y} = \mu_2 f_0 \frac{\partial^2 u_2}{\partial y^2} + k(u_2 - u_1),
\]

\[
\rho_2 u_2 \frac{\partial v_2}{\partial x} + \rho_2 v_2 \frac{\partial v_2}{\partial y} = \mu_2 f_0 \frac{\partial^2 v_2}{\partial y^2} + k(v_2 - v_1),
\]

and in the mentioned above boundary conditions only the condition \(v_2 = f_3(\tau_1)\) will be changed to \(v_2 = f_3(\tau_1, \tau_2)\).

The changed system of equations with changed boundary conditions was solved numerically by means of application of the finite-difference method. According to figure 4, shear stress of the second phase is increasing to a certain value while shear stress of the first phase is decreasing. Therefore, stress which is created by the second phase (erythrocytes) can reach the same order of magnitude as the carrying phase. Large stress inside the second phase (erythrocytes) may potentially lead to the endothelial damage and erythrocyte deformation [16]. To check this statement, experiments were conducted (in accordance to methods discussed above). The results of these experiments are given below.

**Damage at accelerated flow**

Table 1 shows the results of the experiments. According to table 1, experimental damage (3rd column) considerably exceeds the theoretically calculated damage. In addition, as it can be seen from table 2, instead of decreasing, the damage along the BL increases. Theoretical and experimental results are similar up to 0.015. (columns 3 and 5). After this point, damage can no longer be caused by LBL. It is also not caused by the stress created by the second phase, because stress at the second phase is determined by increasing inflow of the destroyed particles. Perhaps it is transferred into a transition zone which is characterized by turbulent spots in the LBL or into the turbulent mode where shear stress is determined by chaotic turbulent flow and leads to significant effective mixing of the flow.

---

**Figure 3**

The distribution of transverse velocity. Along the horizontal axis: dimensionless distance; along the vertical axis: dimensionless air-stream velocity. Legend on the right: phases.

**Figure 4**

Change in shear stress of 1st and 2nd phase along the longitudinal coordinate (x). Horizontal axis: distance (m); vertical axis: shear stress \((\frac{N}{m^2})\). Legend on the right: first and second phases.
Experiments inside the TBL

The results of experiments are given in Table 2. According to Table 2, the relative agreement of the results of damage begins from the distance of more than 0.25 (i.e., in the TBL). On distance from 0.15 m to 0.25 m it is characterized by a transitional zone (turbulent spots in LBL). Physically it is explained by the fact that at BLS, the flow inside the BL is changed significantly. Vortices are formed around the curved surface. Transformed flow interacts with the flat surface of the undamaged wall. These transformed flows also change flow conditions (i.e., transfer into the turbulent mode).

Therefore, short-term compression caused by injury can lead to a transfer of a LBL into a TBL. Let's calculate shear stress in the TBL. We will use the law of resistance of Blasius [7,8] for Reynolds number less than 100,000:

\[ \tau = 0.03955 \frac{u_d^2}{\rho} \left( \frac{U}{\nu} \right)^{0.25} \]

If dog was injured, then \( d = 0.015, \nu = 3.78 \cdot 10^{-6}, \rho = 1.056 \) compression by 25% [7]. Results are given in figure 6.

According to figure 6, when velocity is greater than \( 2.4 \frac{m}{s} \), stress exceeds physiological value and the damage to the endothelium occurs. For human aorta, such values begin from 35% of compression (for instance, compression by 40% leads to turbulent stress of \( 47.52 \frac{N}{m^2} \) at peak velocity of \( 1.9 \frac{m}{s} \)). Therefore, if detrimental compression happens when aortic velocity is greater than \( 2.4 \frac{m}{s} \), might lead to damage of endothelium. (Normal velocities in canine aorta are ranging from \( 0.4 \frac{m}{s} \) to \( 2.9 \frac{m}{s} \).) Let's examine stress appearing on the vessel wall as the result of compression. Results are given in figure 7:

According to figure 7, compression of 10% or greater leads to the damage to the layer of endothelium. Calculations, however, are given at peak velocity \( 2.9 \frac{m}{s} \). Therefore, only certain injuries can lead to the damage of the layer of endothelial cells.

Summary of results

1. Injury that caused sudden compression of the aorta that happened within the range of Reynolds number of 2,000 – 20,000 and within 0.039 to 0.392 s leads to BLS on the curved surface of the vessel.

2. The curved surface of the vessel (upon sudden acceleration with compression caused by injury) results in shear stress that can exceed the physiological value.

3. Damage of the vessel wall is the function of the shear stress.

4. Shear stress of the second phase (2nd phase) can reach the same order of magnitude as the carrying phase (1st phase).

5. Sudden compression caused by injury may lead to transfer of LBL into TBL.

6. If compression of 25% or more and velocity greater than \( 2.4 \frac{m}{s} \) or compression of 10% or more and velocity greater than \( 2.9 \frac{m}{s} \) occur simultaneously, then a critical value is reached, which causes the damage to the layer of endothelium. The high peak velocity and/or significant compression are necessary to cause the damage to the layer of endothelial cells.

Discussion and conclusion

The mathematical model constructed in our study estimates the short-term aortic hemodynamic conditions at the time of chest or abdominal contusion. At present these properties cannot be determined directly in vivo.
Navier-Stokes equations, multiphase equations and BL equations have been used for decades to describe fluid dynamics; however, this study is the first attempt to describe hemodynamic variables at the moment of injury. Traumatic agents cause different types of blood flow changes depending on the location of the trauma. If the chest or pelvis were injured, then the blood flow motion, although moderately accelerated, may not become sudden because the rib cage or iliac wings act as a shock absorber. If, however, there were an abdominal injury or anterior-posterior pelvic fracture, then the acceleration is significant and, therefore, the blood flow motion is sudden.

According to autopsy findings, the majority of patients with aortic injury have either rib and/or sternal fractures indicating significant chest compression [17]. We report in this work that sudden aortic compression caused by trauma leads to separation in the BL. The study of flow separation from the surface of a vessel wall and the determination of changes in the blood flow that develop as a result of the separation are among the most fundamental and difficult problems of hemodynamics. To the best of our knowledge, a clear mathematical explanation of the physical processes leading to the separation has not been reported so far.

Observations of the importance of BLS include its role as the initiating factor in the development of intimal and neointimal hyperplasia at vascular anastomoses [18, 19]. BLS has been shown to predispose patients to atherosclerosis at arterial bifurcations [20]. LoGerfo also suggested that BLS is an important factor in the localization and progression of the atherosclerosis rather than an initiator of the disease [6].

There is evidence that high shear stress may cause endothelial injury [21]. However, the shear stress required to damage the endothelium is far beyond that found in the normal physiological conditions [7]. We found, however, that the value of shear stress is directly related to the value of the force of acceleration that results from injury. Trauma, therefore, due to possible significant compression might cause an increase in the value of shear stress as much as 3-fold. In the study by Ochsner et al. [22] 83% of aortic ruptures occurred in patients with an anterior-posterior pelvic fracture pattern, and an incidence of aortic rupture was more than 9 times greater in these patients than the incidence of aortic rupture in the overall blunt trauma population (7.5% vs. 0.8; p < 0.001). Interestingly, there was no increased incidence of aortic rupture among patients with any other pelvic fracture pattern, such as lateral compression where iliac wings might possibly act as a shock absorber.

### Table 1: Damage at accelerated flow

| Distance (m) | Air-stream velocity (m/s) | Damage (g/s) | Theoretical stress (N/m²) | Theoretical damage (g/s) |
|-------------|---------------------------|--------------|---------------------------|--------------------------|
| 0.09        | 7.46                      | 0.09         | 0.09                      | 0.094                    |
| 0.15        | 8.2                       | 0.098        | 0.09                      | 0.097                    |
| 0.21        | 8.9                       | 0.28         | 0.087                     | 0.094                    |
| 0.25        | 9.4                       | 0.46         | 0.078                     | 0.093                    |
| 0.31        | 10.13                     | 0.83         | 0.086                     | 0.094                    |
| 0.35        | 10.61                     | 1.09         | 0.087                     | 0.094                    |

### Table 2: Damage inside laminar boundary layer

| Distance (m) | Air-stream velocity (m/s) | Theoretical damage (g/s) | Damage (g/s) | Incident flow velocity (m/s) |
|-------------|---------------------------|--------------------------|--------------|------------------------------|
| 0.09        | 7.46                      | 0.09                     | 0.21         | 8.5                         |
| 0.15        | 8.2                       | 0.1                      | 0.58         | 9                            |
| 0.21        | 8.9                       | 0.28                     | 0.657        | 9.5                         |
| 0.25        | 9.4                       | 0.46                     | 0.88         | 9.9                         |
| 0.31        | 10.13                     | 0.83                     | 1.16         | 10.43                       |
| 0.35        | 10.61                     | 1.09                     | *            | laminar flow                 |

* Turbulent flow can not be created at this distance.
We report that sudden compression as the result of the injury leads to blood turbulence which is one of the seven known agents that contribute to human atherogenesis [23,24]. Generally, the presence of turbulence in the cardiovascular system is an indication of some type of abnormality [23,24]. For instance, depending on the phase of the occurrence of turbulence during the cardiac cycle, turbulence near a cardiac valve might indicate either valvular stenosis or regurgitation [23,24]. Turbulence has been shown to be a risk factor for atherosclerotic plaques, since plaques predominate in turbulent flow areas of the bloodstream, such as points of arterial branching [23,24].

Our research has the potential for improvement of safety features of motor vehicles. The design of seatbelts and airbags should take into account the forces acting on large blood vessels as well as associated pressures. The data on injury from crashes can be used to compare injury outcomes from seat belt use or airbag deployment and to predict forces affecting the chest. The collision reports will become the basis for classification of actual forces that are produced by a motor vehicle crash. On the other hand, this work may become a basis for comparison of injury data with physical forces to establish a relationship between the range of injury and the nature of forces involved under different scenarios of deployment of safety features. This information will be useful for maximizing the protective effect of seat belts, airbags and other safety devices.

**Conclusion**

Certain injury conditions can cause damage to the endothelium (for instance, the significant degree of compression resulting from injury). Our research may point up new ways of reducing the damage from blunt trauma to large vessels. It has the potential for improvement of safety features of motor vehicles. This work will better our understanding of the precise mechanics and critical variables involved in diagnosis and prevention of blunt trauma to large vessels.

**Abbreviations**

Aerodynamic tube (AT)

Boundary layer (BL)

Boundary layer separation (BLS)

Laminar boundary layer (LBL)

Turbulent boundary layer (TBL)

**Authors’ contributions**

RI designed the study, assisted in performing experiments, performed mathematical analysis and drafted the manuscript. NS participated in the design of the study, drafted the manuscript and assisted in mathematical analysis. JS and LK conceptualized the research and assisted in drafting the manuscript. IK performed experiments and assisted in mathematical analysis.
References

1. Feczko JD, Lynch L, Pless JE, Clark MA, McClain J, Hawley DA: An autopsy case review of 142 nonpenetrating (blunt) injuries of the aorta. J Trauma 1992, 33:846-9.

2. Parmenter LF, Mattingly TW, Manion WC. Nonpenetrating traumatic injury of the aorta. Circulation 1958, 17:1086-1100.

3. Ahrran K, Smith DC: Trauma to the aorta and aortic arch branches. Curr Opin Cardiol 1998, 13:353-68.

4. Malhotra AK, Fabian TC, Croce MA, Weiman DS, Gavant ML, Pate JW: Minimal aortic injury: a lesion associated with advancing diagnostic techniques. J Trauma 2001, 51:1042-8.

5. Fox RW, McDonald AT, Fox RW, McDonald AT: Introduction to fluid mechanics. New York: John Wiley; 1975.

6. LoGerfo FW, Soncrant T, Teel T, Dewey CF Jr: Boundary layer separation in models of side-to-end arterial anastomoses. Arch Surg 1979, 114(12):1369-73.

7. Caro CG, Pedley TJ, Schroter RC, Seed WA: The mechanics of the circulation. Oxford: Oxford University Press; 1978.

8. Schlichting H: Boundary layer theory. New York: McGraw-Hill Book Co; 1968.

9. Hakaim AG, Nalbandian MN, Heller JK, Chowla AC, Oldenburg WA: Improved patency of prosthetic arteriovenous grafts with an acute anastomotic angle and flow diffuser. J Vasc Surg 2003, 37(5):1032-5.

10. Oury JH, Doty DB, Oswalt JD, Knapp JF, Mackey SK, Duran CM: Cardiopulmonary response to maximal exercise in young athletes following the Ross procedure. Ann Thorac Surg 1998, 66(6):153-4.

11. Rakhmatulin KhA: The basic gas dynamics of interpenetrating motions of compressible media (in Russian). Prikl Mat Mekh 1956, 20:185-195.

12. Navruzov KN, Khakberdiev ZhB: Dynamics of Non-newton fluids. Tashkent: Tashkent 2000.

13. Joris I, Zand T, Majno G: Hydrodynamic injury of the endothelium in acute aortic stenosis. Am J Pathol 1982, 106:394-408.

14. Caro CG, Fitz-Gerald JM, Schroter RC: Atheroma and arterial wall shear. Observation, correlation and proposal of a shear dependent mass transfer mechanism for atherogenesis. Proc R Soc Lond B Biol Sci 1971, 177:109-59.

15. Fauzulleev FD: Laminar motion of multiphase media in conduits. New York: Consultants Bureau; 1969.

16. Sabah HN, Stein PD: Effect of erythrocytic deformability upon turbulent blood flow. Biochemistry 1976, 13(5):309-314.

17. Shkrum MJ, McClafferty KJ, Green RN, Nowak ES, Young JG: Mechanisms of aortic injury in fatalities occurring in motor vehicle collisions. J Forensic Sci 1999, 44:44-56.

18. Imparato AM, Bracco A, Kim GE, Zeff R: Intimal and neointimal fibrous proliferation causing failure of arterial reconstructions. Surgery 1972, 72:1007-17.

19. Oblath RW, Buckley FO Jr, Green RM, Schwartz SI, DeWeese JA: Prevention of platelet aggregation and adherence to prosthetic vascular grafts by aspirin and dipryridamole. Surgery 1978, 84:37-44.

20. Gutstein WH, Schneck DJ, Marks JO: In vitro studies of local blood flow disturbance in a region of separation. J Atheroscler Res 1968, 8:381-8.

21. Fry DL: Acute vascular endothelial changes associated with increased blood velocity gradients. Circ Res 1968, 22:165-97.

22. Osthener MG, Hoffman AP, DiPassqua D, Cole EJ, Rzeczycki GS, Webster DW, Champion HR: Associated aortic rupture-pelvic fracture: an alert for orthopedic and general surgeons. J Trauma 1992, 33(3):229-34.

23. Bluesten D, Emav T: Transition to turbulence in pulsatile flow through heart valves – a modified stability approach. J Biomech Eng 1994, 116:477-87.

24. Constantinides P: The role of arterial wall injury in atherogenesis and arterial thrombogenesis. Zentralbl Allg Pathol 1989, 135(6):517-30.