AN OPTICAL–INFRARED STUDY OF THE YOUNG MULTIPOLAR PLANETARY NEBULA NGC 6644

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ABSTRACT

High-resolution Hubble Space Telescope imaging of the compact planetary nebula NGC 6644 has revealed two pairs of bipolar lobes and a central ring lying close to the plane of the sky. From mid-infrared imaging obtained with the Gemini Telescope, we have found a dust torus which is oriented nearly perpendicular to one pair of the lobes. We suggest that NGC 6644 is a multipolar nebula and construct a three-dimensional model that allows the visualization of the object from different lines of sight. These results suggest that NGC 6644 may have similar intrinsic structures as other multipolar nebulae and the phenomenon of multipolar nebulosity may be more common than previously believed.

Key words: ISM: general – planetary nebulae: individual (NGC 6644) – stars: individual (AGB) Online-only material: color figures

1. INTRODUCTION

Although planetary nebulae (PNs) are usually assumed to have a simple spherical geometry, the actual morphologies of PNs have diverse forms. When deeper, narrow-band observations are taken, PNs often reveal outer structures such as lobes and haloes. While it has been recognized for some time that there are many examples of PNs with prominent bipolar structures (e.g., NGC 2346, NGC 6302), an increasing number of nebulae with multiple lobe structures have been found (e.g., NGC 2440, NGC 6072). The origin of such structures and the physical mechanisms leading to their formation are yet to be understood.

Since PNs are made up of ionized, molecular, and dust components, a comprehensive understanding of the dynamical processes that create the observed morphological structures can only be obtained with multiwavelength observations. With modern 8 m class telescopes equipped with high-resolution mid-infrared imaging capabilities, the distribution of the dust component can be compared to the distribution of the ionized gas component with comparable resolution. In particular, the infrared observations will allow us to search for the existence of a high-density torus and determine what role the torus play in the collimation of observed outflows.

The PN NGC 6644 (PNG 008.3-07.3, Hen 2-408, VV 188, IRAS 18295–2510) was discovered by Hubble (1921). Based on early ground-based imaging, this compact nebula has been classified as an elliptical PN by Stanghellini et al. (1993). It has a large heliocentric radial velocity of 205 km s\(^{-1}\) and is suggested as a PN in the Galactic Bulge (Beaulieu et al. 1999). However, Góny et al. (2004) argued that this object is not a Galactic Bulge PN based on their spectroscopic study. Although NGC 6644 has a relatively high surface brightness, it has received very little attention in the literature. There are only three entries in the Astrophysical Data System on the object after a search of keywords in the abstract, and the papers are mostly related to abundance studies of the nebula. Aller & Keyes (1988) proposed that the object was made from a much less metal-rich mixture than the Sun.

In this paper, we present results of our morphological study for NGC 6644 based on optical and infrared images taken with Hubble Space Telescope (HST) and Gemini Telescope, respectively. These images reveal for the first time unique morphologies and properties of the system. The observations and data reductions are described in Section 2. In Section 3, we present the results of imaging and spectroscopy in the visible and mid-infrared for this nebula. A discussion of the results and a comparison of the observations with a three-dimensional (3D) model are presented in Section 4. Finally, a conclusion is given in Section 5.

2. OBSERVATIONS AND DATA REDUCTION

2.1. HST Narrow-band Imaging

Our HST data on NGC 6644 were retrieved from the Space Telescope Science Archive. The images of NGC 6644 were obtained under program 8345 (PI: R. Sahai) using the Wide Field Planetary Camera 2 (WFPC 2) on HST. The object was observed with the Planetary Camera on 2000 February 26, which provides a 36.8′ × 36.8′ field of view (FOV) at a spatial resolution of 0.045 pixel\(^{-1}\). The actual observations were made with different exposure times (from 20 s to 400 s) to allow for the imaging of both the bright central region and the faint outer parts. The data were processed through the HST pipeline calibration. Standard bias subtraction and flat-field correction were performed. Data were taken in two-step dithered positions to enhance spatial sampling and cosmic rays removal by using the task crej in the STSDAS package of IRAF. The processed F656N (H\(_\alpha\)) observation (λ\(_c\) = 6564 Å, Δλ = 22 Å) with a total exposure time of 1120 s is shown in Figure 1.

2.2. Gemini T-ReCS Observations

The imaging observations were made with the Thermal-Region Camera Spectrograph (T-ReCS) on the Gemini South Telescope under program GS-2007A-DD-9 on 2007 October 10. The FOV of the detector is 28′′ × 21′′ with a pixel scale of 0′′.09. The background emission from the sky and telescope were removed with chopping and nodding during the observations. NGC 6644 was imaged with three narrow/medium-band filters: [S IV] (λ\(_c\) = 10.52 μm, Δλ = 0.17 μm), PAH-2 (λ\(_c\) = 11.30 μm, λ\(_c\) = 10.52 μm, Δλ = 0.17 μm) and PAH-3 (λ\(_c\) = 11.70 μm, Δλ = 0.50 μm).
The mid-infrared spectra of NGC 6644 were obtained by the Infrared Spectrograph (IRS; Houck et al. 2004) through the observation program 3633 (PI: M. Bobrowsky) with Astronomical Observation Request (AOR) key of 11334400. The measurements were observed using the Short-Low (SL) module (5.2–14.5 μm), Short-High (SH) module (9.9–19.6 μm), and the Long-High (LH) module (18.7–37.2 μm) with spectral dispersions of $R \sim 600$. The diaphragm sizes are $3.6 \times 57''$, $4.7 \times 11.1''$, and $11.1 \times 22.3''$ in SL, SH, and LH modules, respectively. The total integration time of IRS observation is 347 s. The IRS measurement of NGC 6644 was not positioned at the central star but on the outer lobe. An overlay of the IRS apertures on the optical image of the nebula is shown in Figure 2.

2.3. Spitzer IRS Spectrum

The mid-infrared spectra of NGC 6644 were obtained by the Infrared Spectrograph (IRS; Houck et al. 2004) through the observation program 3633 (PI: M. Bobrowsky) with Astronomical Observation Request (AOR) key of 11334400. The measurements were observed using the Short-Low (SL) module (5.2–14.5 μm), Short-High (SH) module (9.9–19.6 μm), and the Long-High (LH) module (18.7–37.2 μm) with spectral dispersions of $R \sim 600$. The diaphragm sizes are $3.6 \times 57''$, $4.7 \times 11.1''$, and $11.1 \times 22.3''$ in SL, SH, and LH modules, respectively. The total integration time of IRS observation is 347 s. The IRS measurement of NGC 6644 was not positioned at the central star but on the outer lobe. An overlay of the IRS apertures on the optical image of the nebula is shown in Figure 2.

Data were reduced starting with basic calibrated data from the Spitzer Science Center’s pipeline version s18.7 and were run through the IRSCLEAN program to remove rogue pixels. Next the SMART analysis package (Higdon et al. 2004) was used to extract the spectra. A final spectrum was produced using the combined IRS observations to improve the signal-to-noise ratio (S/N). Since the IRS spectrum in the short wavelength range is taken from the outside of the nebula, some scaling is needed for the shorter wavelength observations. In Figure 3, we scale the IRS SH and SL observations by factors of 1.29 and 1.625, respectively and are able to obtain a smooth spectrum. When we compare these adjusted flux levels to the photometric measurements obtained from the four Gemini T-ReCS filters, we find good agreement (Figure 3). Further confidence of this adjustment is found by the agreement between the spectra and the Infrared Astronomical Satellite (IRAS) 12 and 25 μm fluxes. Also plotted in Figure 3 are the four Gemini filter transmission functions. We can see that the [S\textsc{iv}] and PAH-2 filters have most of their fluxes contributed by the [S\textsc{iv}] line and the Si-5 filter mainly takes in fluxes from the 12 μm plateau emission feature and the $Q_{\alpha}$ filter measures the general dust continuum. Although the $Q_{\alpha}$ filter covers the [S\textsc{iii}] line, we estimate that the contribution of the [S\textsc{iii}] line to the filter flux to be $\sim 8\%$ and should not affect the observed morphology of the dust continuum.

3. RESULTS

3.1. Lobes and Halo

The HST Hα image of NGC 6644 reveals several morphological structures in the nebula (Figure 1). Easily discernible...
are two bipolar lobes (labeled as $a$–$a'$ and $b$–$b'$) along the approximate SE–NW direction. A ring structure (labeled as “ionized torus”) near the center can also be seen. Beyond the bright ring, there are associated nebulosities (marked as $c$), in particular in the western direction. It is possible that these nebulosities represent the projection of a third pair of lobes which are aligned nearly along the line of sight. An extended diffuse structure (marked as “halo”) can also be seen. In Figure 4, we display the image with the central part of the nebula saturated, allowing the fainter outer structures to be seen. There are apparently sub-structures in the nebula. In Figure 5, we mark a few possible sub-structures (which appear predominantly in the north) as $d$, $e$, $f$, $g$, $h$. Closer to the core of NGC 6644 appears

Figure 3. Comparison between the Spitzer IRS spectra and the Gemini T-ReCS photometry. The observations of the IRS SH and SL have been adjusted upward by factors of 1.29 and 1.625, respectively (see the text). The T-ReCS and IRAS photometric measurements are shown as circles and squares, respectively. The Gemini T-ReCS filter transmission functions are plotted as dotted lines.

(A color version of this figure is available in the online journal.)

Figure 4. HST $H\alpha$ image of NGC 6644 with an intensity display setting to show the outer structures of the nebula. The intensity display is on a logarithmic scale and the gray scale bar is given at the bottom in units of counts per pixel.

Figure 5. A number of faint sub-lobes can be seen in the HST $H\alpha$ image of NGC 6644. The caps of the lobes have been highlighted by bars for easy identification. The lines perpendicular to the bars labeled as $d$, $e$, $f$, $g$, $h$ are inferred axes of the lobes.

(A color version of this figure is available in the online journal.)

Figure 6. Three inner lobes are marked on the HST $H\alpha$ image of NGC 6644. The labels are as in Figure 5.

(A color version of this figure is available in the online journal.)

several more sub-substructures, again most obvious in the northern part of the image. These are labeled $i$, $j$, $k$ in Figure 6. The axes of these structures are defined by the caps of the lobes, and locations are marked by horizontal bars in Figures 5 and 6. The orientation of the axis of $d$ is less well defined as the cap is not complete, but the others are reasonably well defined. What is interesting is that the axis $j$ seems to lie along the same orientation as $d$, as well as the axes of $k$ and $g$. The fact that these
Figure 7. Profile of surface brightness distribution of Hα emission in NGC 6644 averaged over all angles except the regions occupied by the lobes. A young PN. The halo can be fitted by a density law of $r^{-2}$ (model in solid line) but the intensities in the inner regions require contribution from a separate component (dashed line). See the text for details.

(A color version of this figure is available in the online journal.)

The symmetry axes of lobes $a$ and $b$ intersect approximately at the position of the central star. The two pairs of butterfly-shaped, closed-end bipolar lobes ($a-a'$ and $b-b'$) have symmetry axes that lie at the position angles (PAs) of $-43^\circ \pm 2^\circ$ and $-27^\circ \pm 2^\circ$, respectively. The lobes $a-a'$ and $b-b'$ have similar projected sizes on the sky. Lobe $a-a'$ has a projected length of $\sim 7.76$ arcsec and $b-b'$ has a projected size of $\sim 8.81$ arcsec. If we assume a distance of 2.5 kpc (Cahn et al. 1992) the physical size of the total extent of the $b-b'$ lobe is $\sim 0.35$ sec $\theta$ light year, where $\theta$ is the inclination angle. For an expansion velocity of 100 km s$^{-1}$, the kinematic age of this lobe is $\sim 521$ years if the lobes lie close to the plane of the sky. From this kinematic age, NGC 6644 can be considered as a young PN.

For the central ring (“ionized torus”), its major axis is found to have an angular orientation of PA $\sim -55^\circ \pm 3^\circ$. Similar structures have been found in other multipolar PNs such as the two bright partial rings seen in the waist of He 2-47 (Sahai 2000b, his Figure 1(b)). The eastern segment of the ring in NGC 6644 is brighter than the western side, as evident in both the Hα image (Figure 1) and in the [S iv] image in Figure 8. This suggests that the eastern segment is closer to us, whereas the western side suffers from extinction from intervening dust in the dust torus.

This bright multipolar nebula is surrounded by a faint, extensive halo with a well-defined spherical shape. The extended halo emission at Hα can be seen in an averaged radial surface brightness profile (with the two bipolar lobes excluded) shown in Figure 7. This profile is constructed from measurements of $5^\circ$ intervals between PA $= 25^\circ - 85^\circ$ and PA $= 205^\circ - 265^\circ$, after the removal of all field stars. The intensity of the nebula within $1.4'$ is dominated by the ionized torus and the angular radius of the faint halo is estimated to be larger than $4\arcsec$. The surface brightness of the extended part at radius $4\arcsec$ relative to that in the peak is $\sim 10^{-3}$. At a distance of 2.5 kpc, the physical size of the halo is 0.15 light year, with a kinematic age of $\sim 3 \times 10^3$ yr, if the expansion velocity of the halo is 15 km s$^{-1}$.

An intensity profile of the nebula (minus the lobes) averaged over all angles is shown in Figure 7. Also shown is a fit of the profile by a halo density distribution of $r^{-2}$, where $r$ is the radial distance from the central star, with an inner and outer radii of $R_{\text{in, out}} = 0.585$ and 6.5 arcsec, respectively. For the modeling, we have assumed that Hα is optically thin and its intensity integrated along each line of sight is proportional to the emission measure, $EM(p) = \int p n^2(r) d\ell$, where $n(r)$ is the density of hydrogen ion, $p$ is the angular distance from the central star, and $d\ell$ is the path length element along the line of sight. Therefore, under the assumption that the halo is a spherical shell we have

$$EM(p) = 2n^2(R_{\text{in}})R_{\text{out}}^{-2a}p^{1-2a} \int_{\theta_1}^{\theta_2} (\cos \theta)^{2a-2} d\theta, \quad (1)$$

where $\theta_1 = \cos^{-1}(p/R_{\text{in}})$ and 0 for $R_{\text{in}} \leq p \leq R_{\text{out}}$ and $p < R_{\text{in}}$, respectively, and $\theta_2 = \cos^{-1}(p/R_{\text{out}})$. For a stellar wind with a constant mass-loss rate and expansion velocity, $a$ has a value of 2 and the density distribution varies as $n(r) = (R_{\text{in}}/r)^{-2}$. While the model provides a reasonably good fit to the halo, it fails in the inner region.

In a second model, we introduce a face-on torus with a cross-section radius of $R_t = 0.4'$ and a homogeneous density distribution of $n_t$. For the torus, we have

$$EM(p) = 2n_t^2 \sqrt{R_t^2 - (p - R_t)^2}, \quad (2)$$

where $R_c = 0.585$ is the distance from the circle center of the torus to the central star. The torus is embedded in a halo with a inner radius of 1 arcsec. This addition makes an improved fit to the intensity profile in the central region. The remaining excess in emission between 1 and 2 arcsec (as seen in Figure 7) could be due to emission from the nearby pole-on lobes.

In the model, we do not consider the effect of extinction. Due to the presence of dust torus (see below), the central region may have a higher extinction than the halo. If this is the case, $EM(p)$ given in Equation (2) could be underestimated and the actual intensity of the inner region would be even higher. This will further strengthen the need for a face-on torus.

3.2. The Dust Torus

In Figure 8, we show the four-band T-ReCS (PAH-2, Si-5, [S iv], and Qλ) flux-calibrated images in units of Jy per square arcsecond. The T-ReCS images of NGC 6644 were each deconvolved using the observations of the standard star HD 175775 taken immediately after the science observations. The “lucy” task in the “analysis” group of the “stsci” external IRAF package was used for the deconvolutions. For each filter a background-subtracted, normalized point-spread function (PSF) image of size 51 by 51 pixels was created from the standard star observation and this was used as the kernel for the deconvolution. The deconvolution was iterated 25 times to produce the final deconvolved images. It was found that most of the improvement in the images took place within the first 5 to 10 iterations. Under the seeing conditions reported in Section 2.2, the Lucy deconvolution is robust to small changes in the FWHM of the PSF. The deconvolved images are shown in Figure 9 together with the Hα image on the same scale.

From Figures 8 and 9, we can clearly see that there is a spatially resolved structure in the central region of NGC 6644.
Figure 8. Flux-calibrated Gemini T-ReCS images of NGC 6644. The four panels give, respectively, the PAH-2 11.3 μm image (upper left), the Si5 11.66 μm image (upper right), [S iv] 10.52 μm narrow-band image (lower left), and Qa 18.3 μm medium-band image (lower right). North is up and east is to the left. The intensity display is on a linear scale and the gray scale bar is given at the bottom in units of counts per pixel.

The 10.5 μm [S iv] image shows a partial ring structure similar to that seen in the central region in the Hα image, suggesting that they both trace the ionized gas distribution. The Si5 and PAH2 images are different from the [S iv] image, showing a double-peaked structure. The Qa image shows a more diffuse structure than the other three images. The overall maximum angular extent of the emission regions in the four images is similar. If the double-peaked structure in the PAH2 and Si5 filter images is interpreted as an oblique torus, the size of the torus has dimensions of 28 × 15 pixels with the major axis lying along PA ~ −44° ± 3°. This translates to a major axis size of 2.52 and a minor axis size of 1.35. Assuming that this ellipse is a projection of a circle on the sky, an angle of tilt of the dust torus is approximately 32° (with 90° being pole on). The fact that this torus is not obvious in the optical image is probably because the dust torus is optically thick to the UV photons and the volume of the torus is not ionized.

A more detailed comparison between the HST Hα image with the Gemini [S iv] image shows that the ring in the [S iv] image is less complete than its counterpart in the Hα image (Figure 9).
Figure 9. Comparison between the HST Hα image (lower left) and the flux-calibrated, deconvolved, false-color Gemini T-ReCS images displayed on the same pixel scale. The Gemini [S\textsubscript{iv}] 10.52 \textmu m narrow-band, PAH-2 11.3 \textmu m narrow-band, and the Q\textsubscript{a} 18.3 \textmu m medium-band images are given in the upper left, upper right, and lower right panels, respectively. (A color version of this figure is available in the online journal.)

Table 2

| Filter Name | Central Wavelength $\lambda_c$ (\textmu m) | Bandwidth $\Delta \lambda$ (\textmu m) | Flux Density (Jy) | IRS Flux Density$^a$ (Jy) |
|-------------|---------------------------------|---------------------------------|-----------------|-----------------|
| [S\textsubscript{iv}]| 10.52                        | 0.17                            | 4.18            | 4.29            |
| PAH2        | 11.30                          | 0.61                            | 2.31            | 2.41            |
| Si5         | 11.66                          | 1.13                            | 2.39            | 2.32            |
| Q\textsubscript{a} | 18.30                        | 1.38                            | 5.78            | 5.01            |

Note. $^a$ The IRS flux density is obtained from integrating the convolution of filter transmission curve and Spitzer IRS spectrum.

There are two factors that can cause a difference between the two images. The [S\textsubscript{iv}] image includes contributions from both the [S\textsubscript{iv}] line and dust continuum and therefore should reflect the morphology of the ionized torus as well as the oblique dust torus. Since the dust torus (see Figure 8) is brighter in the eastern side, the [S\textsubscript{iv}] image should therefore appear as a less complete ring than in the Hα image. If there is also effect of local extinction that causes the eastern side to be brighter (see Section 3.1), then the [S\textsubscript{iv}] image should suffer from less extinction and therefore should show a more complete ring. This effect is apparently not as important.

In Table 2, we compare the fluxes observed in the four-band T-ReCS observations and fluxes from the IRS observations. Columns 1 and 2 list the observed filter names and central wavelengths of these filters. The corresponding bandwidth for each filter is given in Column 3. The measured integrated fluxes from the T-ReCS observation are given in Column 4. The simulated T-ReCS in-band fluxes derived from the IRS spectrum and the filter profiles of the four bands are given in Column 5. There is general consistency between the Gemini and IRS measurements with the discrepancies being due to varying sky conditions and uncertainty of photometric flux calibrations. These uncertainties are higher in the Q\textsubscript{a} band.

3.3. Infrared Features

The combined IRS spectrum of NGC 6644 shown in (Figure 3) shows a strong infrared continuum due to dust emission peaking at $\sim$30 \textmu m. On top of this continuum are a number of emission lines typical of the spectrum of PNs. Among the strongest emission lines detected are the fine structure lines of [S\textsubscript{iv}] at 10.51 \textmu m, [O\textsubscript{iv}] at 25.88 \textmu m, [Ne\textsubscript{iii}] at 15.55 and 36.01 \textmu m, and weaker H recombination lines H\textsubscript{i} (7–6) at 12.37 \textmu m and H\textsubscript{i} (8–7) at 19.07 \textmu m (Figure 10). The measured emission line fluxes in the spectrum are given in Table 3.
first two columns in Table 3 contain the central wavelengths of the emissions and line identifications, respectively. Column 3 gives the observed fluxes measured using the Gaussian fitting routine. Using the fine-structure lines [S\textsc{iii}] and [Ne\textsc{iii}], we can derive the electron density. From the [S\textsc{iii}] 18.7 \mu m to 33.4 \mu m and [Ne\textsc{iii}] 15.6 \mu m to 36 \mu m line ratios, we derive log \( n_e = 3.8 \) and 3.4 cm\(^{-3}\), respectively. The last column gives the total flux of the entire nebula by scaling the observed flux.
The electron temperature of Te is based on Case B recombination line theory. Assuming an of ne by a factor, derived from the average theoretical flux ratio between infrared H\textsc{i} lines (H\textsc{i} (7–6), H\textsc{i} (8–7)) and H\textsc{\beta} emission based on Case B recombination line theory. Assuming an electron temperature of $T_e = 12,000$ K and an electron density of $n_e = 7,200$ cm$^{-3}$ (Shaw & Kaler 1989), the average theoretical ratios are then combined to derive the scaling factor of 1.01 for the IRS spectrum.

Also present is the strong, broad AIB at 11.3 $\mu$m. This feature is due to the C-H out-of-plane bending of aromatic compounds and is commonly seen in carbon-rich PNs and proto-PNs (Kwok et al. 2001), can also be seen. A recent study of AIB features (Kwok et al. 2007, p. 171) at an electron temperature of 12,000 K and an electron density $n_e = 7,200$ cm$^{-3}$ (Shaw & Kaler 1989), the average theoretical ratios are then combined to derive the scaling factor of 1.01 for the IRS spectrum.

### 3.4. Spectral Energy Distribution

Although PNs were first known for their optical emission-line characteristics, it was realized after the IRAS mission that a significant amount of the total energy output is emitted in the infrared region due to thermal emission by dust grains. A systematic investigation for spectral energy distributions (SEDs) of young PNs covering the wavelength range from ultraviolet (UV) to far-infrared was performed by Zhang & Kwok (1991). Depending on the stage of evolution, it is found that the photospheric, ionized gas, and dust components are the main contributors to the total observed flux. In constructing the SED of NGC 6644, we have made use of data in the astronomical data archives of both ground-based and space-based observations.

In the UV region, we have used observations from the International Ultraviolet Explorer (IUE) low dispersion spectrophotograph. All IUE spectra of the object are produced at the IUE Data Analysis Center (IUEDAC). Reduced data of this nebula are taken in co-added positions of the short wavelength prime (SWP) and long wavelength prime (LWP) spectra to enhance the S/N, respectively. The parts of this spectrum in SWP and LWP are then combined at 1975 Å to produce a single, resultant spectrum.

A journal of the IUE data and Spitzer IRS AOR key of the nebula are given in Table 4. In the infrared, the Spitzer IRS observations are used together with the photometric observations of IRAS. The photometric measurements of NGC 6644 at the four wavelength bands of 12, 25, 60, and 100 $\mu$m are taken from Tajitsu & Tamura (1998).

The $B$ and $V$ magnitudes of the central star of NGC 6644 are observed by Shaw & Kaler (1989). Near-infrared magnitudes of $I, J, H,$ and $K_s$ observed by the Deep Near-Infrared Southern Sky Survey (DENIS) and the Two Micron All Sky Survey (2MASS) are derived from DENIS database and Ramos-Larios & Phillips (2005). A summary of these archival data is given in Table 5.

After correcting the UV spectrum and optical measurements by an extinction value of 0.41 (Shaw & Kaler 1989; which accounts for both circumstellar and interstellar extinction), we fitted the emerging flux ($F_\lambda$ (total)) by a three-component model including a hot central star, a warm gaseous nebula, and a cool dust shell (Zhang & Kwok 1991). The hot stellar component is assumed to be a blackbody of temperature $T_\ast$. The gaseous nebular continuum is the sum of the free–free (f–f), bound–free (b–f), and two-photon emissions using the emission coefficients given in Kwok (2007, p. 171) at an electron temperature $T_e = 12,000$ K and an electron density $n_e = 7200$ cm$^{-3}$ (Shaw & Kaler 1989). The emerging flux is therefore given by

$$F_\lambda (\text{total}) = F_\lambda (s) + F_\lambda (g) + F_\lambda (d),$$  

### Table 3

| Wavelength ($\mu$m) | Identification | Observed Flux ($10^{-10}$ erg cm$^{-2}$ s$^{-1}$) | Normalized Flux ($10^{-10}$ erg cm$^{-2}$ s$^{-1}$) |
|---------------------|----------------|-----------------------------------------------|-----------------------------------------------|
| 10.51               | [S\textsc{iv}] | 0.1098                                        | 0.1109                                        |
| 11.20               | AIB            | 0.1513                                        | 0.1530                                        |
| 11.76               | [Cl\textsc{v}] | 0.0012                                        | 0.0013                                        |
| 12.37               | H\textsc{i} (7–6) | 0.0024                                        | 0.0025                                        |
| 12.81               | [Ne\textsc{iv}] | 0.0055                                        | 0.0056                                        |
| 13.09               | [Ar\textsc{v}] | 0.0019                                        | 0.0020                                        |
| 14.32               | [Ne\textsc{v}] | 0.0225                                        | 0.0227                                        |
| 15.55               | [Ne\textsc{m}] | 0.2366                                        | 0.2389                                        |
| 18.71               | [S\textsc{m}] | 0.0157                                        | 0.0159                                        |
| 19.07               | H\textsc{i} (8–7) | 0.0016                                        | 0.0017                                        |
| 24.32               | [Ne\textsc{v}] | 0.0087                                        | 0.0088                                        |
| 25.88               | [O\textsc{v}] | 0.1440                                        | 0.1454                                        |
| 33.47               | [S\textsc{m}] | 0.0038                                        | 0.0039                                        |
| 36.01               | [Ne\textsc{m}] | 0.0198                                        | 0.0200                                        |

### Table 4

| PN          | Name       | Instrument   | Exposures |
|-------------|------------|--------------|-----------|
| G 008.3-07.3| NGC 6644   | SWP 31711    | 1800      |
|             |            | SWP 01734    | 1800      |
|             |            | LWR 01630    | 2400      |

### Table 5

| Filters | Flux/Flux Density |
|---------|-------------------|
| $B$ (mag)$^a$ | 16.6 |
| $V$ (mag)$^a$ | 15.63 |
| $-\log F (H\beta) (\text{erg cm}^{-2} \text{s}^{-1})^c$ | 11.01 |

**Notes.**

$^a$ From Shaw & Kaler (1989).

$^b$ From DENIS database.

$^c$ From Ramos-Larios & Phillips (2005).

$^d$ From Tajitsu & Tamura (1998), the color- and diameter-corrected IRAS fluxes were given. Note that the flux measured from 100 $\mu$m is an upper limit.
where $F_1(s)$, $F_2(g)$, and $F_3(d)$ are the flux densities of the photospheric continuum of the central star, the nebular continuum, and the dust thermal emission continuum, respectively. The stellar flux is then given by

$$F_\lambda(s) = (\pi \theta^2_\lambda) B_\lambda(T_\star),$$

where $\theta_\lambda$, $T_\star$ are, respectively, the angular radii and the effective temperatures of the central star, and $B_\lambda(T_\star)$ is the Planck functions for temperatures $T_\star$. Given the fact that the B and V photometry lie in the Rayleigh–Jeans side of the blackbody, the value of $T_\star$ cannot be determined precisely. Our best estimate for $T_\star$ is 119,000 ± 8500 K, which is consistent with the previously reported values of 106,000 K (Shaw & Kaler 1989) and 115,000 K (Zhang & Kwok 1991). With an assumed distance of 2.5 kpc (Cahn et al. 1992), the derived luminosity of the central star is $\sim 4600 L_\odot$. According to the model of Schönberner (1981), the above temperature and luminosity of NGC 6644 imply an evolutionary age >6000 years. The value is at least two times higher than the observed kinematic age (Section 3.1). This may suggest that the visible nebulosity (from which we estimate the kinematic age) originates from very late stellar winds, and observations with higher sensitivity are required to detect the faint outer old structures.

The near-IR photometry points suggest that there is excess emission between 2 and 5 $\mu$m which could be due to a cool companion, a hot dust component, or scattering from an unseen disk. We have therefore added a blackbody of temperature 2950 K to the model fit. Using the same assumed distance, this component would have a luminosity of 14 $L_\odot$.

We find that the observed dust emission component is too broad to be fitted by a single blackbody. We therefore fit the dust component by blackbodies of two different temperatures, a warm ($T_{wd}$) and a cold ($T_{cd}$) dust components. The total dust thermal emission is therefore given by the sum of two terms:

$$F_\lambda(d) = F_\lambda(wd) + F_\lambda(cd) = \frac{3M_{wd}Q_1 B_\lambda(T_{wd})}{4\alpha\rho_s D^2} + \frac{3M_{cd}Q_1 B_\lambda(T_{cd})}{4\alpha\rho_s D^2},$$

where $M_{wd}$ and $M_{cd}$ are the masses of warm and cold dust components, $\alpha$ is the grain radius which depends on its physical shape, $\rho_s$ is the density of grain, $Q_1 = Q_0 (\lambda/\lambda_0)^{-\alpha}$ is the grain emissivity function, and $D$ is the distance to the nebula (Kwok 2007, p. 312). According to our fit, the temperatures of the two dust components are 292 K and 123 K. Assuming $Q (1 \mu m) = 0.1$, $\rho_s = 1$ g cm$^{-3}$, $\alpha = 0.1 \mu m$, $\alpha = 1$, and $D = 2.5$ kpc, we obtain $8 \times 10^{-8} M_\odot$ and $3.12 \times 10^{-5} M_\odot$ for the mass of the warm and cold dust components, respectively. Without spectral coverage between 3 to 10 $\mu$m, it is difficult to constrain the temperature of the warm dust component precisely and the above mass estimate for the warm component is therefore subject to a large margin of error.

The SED of the nebula corrected for extinction is shown in Figure 11. We can see that the model gives a reasonable fit to the observed data from UV through the far-infrared. From the total observed nebular flux of $5.3 \times 10^{-10}$ erg cm$^{-2}$ s$^{-1}$, an EM of $2.5 \times 10^{59}$ cm$^{-3}$ was derived. At a distance of 2.5 kpc and an electron density of 7200 cm$^{-3}$, the total ionized gas mass is derived as 0.03 $M_\odot$. From the model fitting, the fraction of total fluxes between 0.1 and 100 $\mu$m from the central stellar, nebular gaseous continuum emission, and dust components are 38%, 18%, and 44%, respectively. The stellar flux shortward of 0.1 $\mu$m is assumed to have been absorbed by the gas component.

From the SED, we can see that there is a broad 30 $\mu$m feature in NGC 6644. This unidentified emission feature is commonly
seen in carbon stars (Volk et al. 2000), post-asymptotic giant branch stars (Szczerba et al. 1999; Hrivnak et al. 2000; Volk et al. 2002) and young PNs.

4. SIMULATED 3D STRUCTURE OF NGC 6644

The observations of NGC 6644 as reported in Section 3.1 suggest that NGC 6644 is a new member in the club of multipolar nebulae. The class of multipolar nebulae, first described by Manchado et al. (1996) and Sahai & Trauger (1998), now includes many prominent members including NGC 2440 (López et al. 1998), He 2–47 and M1-37 (Sahai 2000b), He 2–113 (Sahai et al. 2000), NGC 6881 (Kwok & Su 2005), NGC 6072 (Kwok et al. 2010), etc. Multipolar nebulae, having a point-symmetric rather than axial symmetric structure, need 3D modeling to visualize their structure. With this aim, we have constructed a schematic model of NGC 6644 assuming three pairs of bipolar lobes and an equatorial dust torus using the software program SHAPE (Steffen & López 2006).

SHAPE is a morpho-kinematic modeling tool intended for the analysis of the 3D geometry and kinematic structure of gaseous nebulae. SHAPE uses both kinematic and spatial observations as a guide for the interactive reproduction of model nebulae. Once a 3D model has been constructed, parameters such as PA and inclination can be manipulated in order to study the object from different orientations. In this work, we are interested in studying the complex 3D morphology of multipolar PN for which SHAPE is an ideal tool. We do not attempt to determine physical parameters such as the temperature and density structure and thus we do not perform any radiative transfer or hydrodynamic simulations. Instead, the brightness and location of each component is assigned based on the qualitative analysis of the observed images.

In addition to the observed double bipolar lobes $a$ and $b$, we hypothesize an additional pair of lobes that is perpendicular to the observed ionized torus in Figure 1. This pair of lobes will be referred to as lobe $c$. We further assume that all three pairs of lobes are equal in length and that the $b$ lobes are oriented in the plane of the sky (inclination 0°). With these assumptions, we can then derive an inclination of $31°$ and $−85°$ for the $a$ and $c$ lobes, respectively. The estimated symmetry axes of these three lobes are along PA = $−42°$, $−28°$, and $76°$ for the $a$, $b$, and $c$ lobes, respectively.

We created a dense torus as the counterpart of the infrared torus seen in the Gemini S15 and PAH2 images with a center that coincides with that of the three lobes. If we assume that the observed elliptical torus is a tilted circular disk, then the model's major-to-minor axis ratio implies an inclination of $35°$ (with $90°$ being pole on). This value is in reasonable agreement with our previous estimate of $32°$ based on the ellipticity of the torus (see Section 3.2). In this model, the symmetry axis of lobe $a\rightarrow a'$ is perpendicular to the infrared dust torus. A counterpart of the ring seen in the HST Hα image and the Gemini [S iv] image is also included. The model parameters of these structures are listed in Table 6 and a schematic mesh model is shown in Figure 12.

This simple model can provide insight into what multipolar PNs look like at different orientations. In Figure 13, we present a 3D representation of NGC 6644 viewed from different angles. The bottom left image is the model as seen from Earth, and the orientations of the other images are as labeled. The $x$- and $y$-axes indicated are those of the image plane and the values represent rotations around those axes in degrees.

Figure 12. Comparison of optical–infrared observations and the 3D model of NGC 6644. (a) The gray frame shows the HST Hα image, as in Figure 1. Linear contour plot of the T-ReCS PAH-2 11.3 μm image was overlaid on the HST image (red lines). (b) SHAPE 3D mesh model. The bipolar lobes $a$, $b$, and $c$ are displayed in red, blue, and green, respectively. The equatorial ring observed in the HST Hα image is shown in light blue and the infrared torus observed in Gemini PAH2 image is in white.

(A color version of this figure is available in the online journal.)

When $y = 90°$, the nebula reveals three distinct pairs of lobes with an elliptical disk. This form is similar to that found in two young PNs, He 2–47 and M 1–37 (Sahai 2000b). When $y = 0°$ the bipolar lobes appear closer together and overlap to varying extents. This type of structure is seen, for example, in NGC 6072 (Kwok et al. 2010) and the Frosty Leo Nebula (Sahai 2000a). The present model illustrates the possible complex 3D structures that may be present in other objects and indeed it would be difficult to visualize the structure of NGC 6644 without the aid of such a model. The main purpose of this simulation is to show how complex the structure of a multipolar PN can be and not as a quantitative fit to the observations. A proper model
of a complex object such as NGC 6644 will require kinematic information from all of the morphological features, which can be provided by integral field spectroscopy with large optical telescopes.

5. CONCLUSIONS

Although NGC 6644 appears to be a typical elliptical nebula, high angular resolution and dynamical range optical observations reveal that it in fact has an extremely complicated structure. From HST and Gemini observations, we have identified at least two pairs of bipolar lobes and an infrared dust torus. The orientation of the dust torus is approximately perpendicular to one pair of the lobes. In addition, there is a ring of ionized gas lying almost on the plane of the sky. We suggest that there could be a third pair of bipolar lobes whose axis is perpendicular to this ring. Several other faint sublobes at different radial distances and angular directions can also be seen.

It is clear that NGC 6644 is a multipolar nebula. In order to illustrate its possible 3D structure, we have constructed a model to show what it may look like when viewed from different perspectives. Its simulated rotated images resemble some other multipolar nebulae, suggesting that nebulae with different appearances may have similar intrinsic structures.

The emergence of the class of multipolar nebulae has greatly altered our perception of the morphological structures of PNs. The commonly assumed simple structures of PNs are probably the result of inadequate sensitivity and spatial resolution imaging. Deep imaging of PNs is needed to reveal the true intrinsic structures of PNs.

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