Highly Resolved and Robust Dynamic X-Ray Imaging Using Perovskite Glass-Ceramic Scintillator with Reduced Light Scattering

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All-inorganic perovskite quantum dots (QDs) CsPbX$_3$ (X = Cl, Br, and I) have recently emerged as a new promising class of X-ray scintillators. However, the instability of perovskite QDs and the strong optical scattering of the thick opaque QD scintillator film impede it to realize high-quality and robust X-ray imaging. Herein, the europium (Eu) doped CsPbBr$_3$ QDs are in situ grown inside transparent amorphous matrix to form glass-ceramic (GC) scintillator with glass phase serving as both matrix and encapsulation for the perovskite QD scintillators. The small amount of Eu dopant optimizes the crystallization of CsPbBr$_3$ QDs and makes their distribution more uniform in the glass matrix, which can significantly reduce the light scattering and also enhance the photoluminescence emission of CsPbBr$_3$ QDs. As a result, a remarkably high spatial resolution of 15.0 lp mm$^{-1}$ is realized thanks to the reduced light scattering, which is so far a record resolution for perovskite scintillator based X-ray imaging, and the scintillation stability is also significantly improved compared to the bare perovskite QD scintillators. Those results provide an effective platform particularly for the emerging perovskite nanocrystal scintillators to reduce light scattering and improve radiation hardness.

X-ray has been widely used in probing the inside information of condensed matter subjects non-destructively, enabling broad applications in safety inspection, medical radiography, defect inspection, etc.

1) direct conversion of X-ray photons into electrical signal; 2) indirect conversion by using scintillators to convert X-ray photons to visible photons first and then be detected by photodiode. In the past few years, perovskite semiconductors have demonstrated great potentials in direct X-ray detection due to their exceptional properties such as large X-ray attenuation coefficient, large mobility-lifetime product ($\mu\tau$), and low-cost solution process.$^{[10-18]}$ Recently, all-inorganic CsPbX$_3$ (X = Cl, Br, I) perovskite scintillators have been explored due to their large X-ray absorption efficiency, intense radio-luminescence (RL), low X-ray detection limit, and fast light decay.$^{[1-4,17-19]}$ Which make them highly competitive to the commercial scintillators, for example, CsI:Tl$^{[20]}$ and LuAG:Ce based scintillators.$^{[21]}$ In addition, perovskite-like metal halides featuring self-trapped exciton emissions have revealed good X-ray scintillation with large Stokes shift and high light yield.$^{[22,23]}$ One advantage of inorganic perovskite as scintillator is that there is no need to grow as big crystals as many conventional scintillator to obtain good scintillation yield, in contrast, the scintillation of CsPbBr$_3$ nanocrystal is several orders of magnitudes stronger than CsPbBr$_3$ single crystal.$^{[1]}$ However,
this unique feature also brings dark sides. First, it requires the perovskite nanocrystal film to be at least several hundred micrometers to ensure sufficient X-ray absorption, such thick film is opaque and the light scattering severely affects the imaging resolution. Second, the stability issue could become even more severe when perovskite crystals shrink to nanometer size due to the massively increased surface/bulk ratio. In addition, the water soluble lead of perovskite nanocrystal is always a risk to human and environment. [13,24] Growing QDs inside inorganic glass has been proven as an effective strategy to prevent the degradation and enhance the environmental stability of II–VI and IV–VI QDs, as well as CsPbX3 QDs showing applications in white LED, laser,[28–31] and 3D laser printing.[32] Here, to resolve those above concerns, we in situ grew inorganic perovskite QDs inside transparent and robust matrix as a glass-ceramic (GC) scintillator, which is inexpensive to fabricate and can be easily scaled up. The incorporation of small amount of Eu improves the crystallinity of CsPbBr3 QDs and alters its size distribution, which further enhances the scintillator performance in terms of spatial resolution and light yield. Finally, our CsPbBr3:Eu GC scintillator with significantly reduced light scattering enables X-ray images with high spatial resolution of 15 lp mm−1, even superior to commercial CsI(Tl) scintillator and delivering impressive stability under continuous X-ray illumination and thermal stressing in humid air.

The CsPbBr3 QDs grown in glass matrix was prepared from the transparent precursor glass (PG) by melt-quenching method, whose XRD pattern of PG (Figure 1a) showed diffuse hump without any diffraction peaks due to its amorphous character. Subsequently the PG was annealed at 500 °C and the CsPbBr3 QDs were precipitated inside it, as evidenced by the XRD peaks (Figure 1a) that are well indexed to pure CsPbBr3 cubic phase (JCPDS No. 54–0752). TEM bright filed image of CsPbBr3 GC (Figure 1b) presents dark spherical CsPbBr3 nanocrystals embedded in glass with average diameter of ≈11 nm. And high-resolution TEM image (Figure 1c) shows clear lattice fringe with interplanar distance between two adjacent crystal planes of ≈0.348 nm, in agreement with the (110) plane of cubic CsPbBr3 crystal. Interestingly, the nanocrystals shrink to the typical size of quantum dots (≈5 nm) with the incorporation of Eu, and the size distribution is more uniform shown in Figure 1d. We reason that the doped Eu3+ likely acts as a nucleation agent which can promote rapid growth and precipitation of CsPbBr3, thus decrease the crystal size and improve the uniformity. In order to confirm this, we monitored the in situ growth of CsPbBr3 QDs induced by high-energy electron beam inside PG as shown in Figures S1 and S2, Supporting Information. The TEM images demonstrate the CsPbBr3 QDs are crystallized more rapidly, easily, and uniformly inside Eu doped PG. Since the thickness of the GC is about 2 mm, the perovskite loading among glass is determined to be ≈116 μm based on molar ratio of source materials and their corresponding density, which is similar to the reference nanocrystal film of ≈100 μm thickness. Since Rayleigh scattering intensity is proportional to d6 (diameter of the dots), the light scattering for the Eu doped CsPbBr3 QDs GC, with optimized doping ratio of 1.5% (Figure 1e) can be significantly eliminated, which ultimately translates to reduced signal crosstalk in photodiode array and therefore improve the image resolution. Figure 1f presents the photographs of CsPbBr3 GC with and without Eu3+ dopant as well as conventional CsPbBr3 nanocrystal film, demonstrating distinctly improved transparency for CsPbBr3:Eu GC visually, mainly because of the reduced scattering rather than the difference of absorption. We applied Eu doping strategy to grow other perovskite QDs (CsPbCl3 and CsPbI3) in the same glass matrix and found the sizes of perovskite QDs also uniformly distributed in the matrix as shown in Figure S1, Supporting Information, which demonstrates the effectiveness of this approach.

X-ray photoelectron spectroscopy (XPS) measurements were carried out to study the doping effect of Eu3+ ions. The binding energy of Pb 4f7/2 shifts from 139.2 eV for CsPbBr3 GC to 138.73 eV for CsPbBr3:Eu GC shown in Figure 1g, which may originate from the increased electron density around Pb2+ after Eu3+ doping, since stronger bonding interaction between Eu3+ and Br− may result in less electron donation from Pb2+ to Br− and thereby increase electron density around Pb2+. This result indicates Eu3+ ions occupy the Pb2+ sites in the CsPbBr3 lattice during the CsPbBr3 crystallization under high temperature, since the ionic radius of Pb2+ and Eu3+ are comparable (Pb2+: 119 pm, Eu3+: 95 pm). The XPS spectrum of Eu3+ 3d5/2 shown in Figure 1h with binding energy at about 1135 eV is almost symmetric and can be fitted perfectly with one gaussian-component suggesting only one chemical environment of Eu3+ existing either in glass matrix or in the Pb2+ lattice sites of CsPbBr3 QDs. Based on the above analysis, we reason that the lightly doped Eu3+ (1.5%) ions are almost entirely incorporated into the CsPbBr3 lattice. Due to the unequal charge of Pb2+ and Eu3+, the Eu3+ is also generated in GC as evidenced by the Eu2+ 3d5/2 spectrum with binding energy at 1125 eV (Figure 1i) in order to keep the charge neutrality.[39,40] This phenomenon can be explained by the charge compensation model based on substitution defect mechanism[40] giving rise to reduction of Eu3+ to Eu2+. Figure 1i shows photoluminescence (PL) spectra of CsPbBr3:xEu GC as a function of Eu concentration. The PL intensity of CsPbBr3 increases gradually as the increase of Eu content because of the promoted nucleation of CsPbBr3 QD leading to more specially-confined dots with smaller sizes,[41,42] also reflected by TEM observation in Figure 1b–d delivering 65% PL yield for CsPbBr3:1.5%Eu GC. However, the characteristic Eu3+ emission (inset in Figure 1i) is almost suppressed but still observable. This is in contrast with some previous reports of Eu doped CsPbX3 nanocrystal or glass ceramic, in which the emission intensity of Eu is comparable with or even stronger than that of CsPbX3.[12,13,18,41–45] This discrepancy is likely attributed to our excellent crystallization of CsPbBr3 promoted by Eu dopant, which makes the band-to-band emission of CsPbBr3 extremely strong, and thereby the Eu intensity seems barely observable. The transparent matrix also serves as inherent protection for perovskite quantum dots. As shown in Video S1, Supporting Information, the PL emission of CsPbBr3 nanocrystal film is immediately quenched when dropped into boiling water, while our CsPbBr3:Eu GC is barely affected, proving its toughness under extreme conditions. Besides, the toxic lead is also kept from dissolving into water with the encapsulation of glass matrix, as confirmed by the inductively coupled plasma atomic emission spectrometry (ICP-AES) measurement. (Table S1, Supporting Information). It showed that only 0.093 mg L−1 of Pb2+ was traced in water after the CsPbBr3:Eu GC was kept in water for 1 h, while 2.06 mg L−1 of Pb2+ was detected for the
Figure 1. Fabrication and characterization of CsPbBr$_3$ quantum dots (QDs) inside transparent amorphous matrix. a) XRD patterns of precursor glass (PG), CsPbBr$_3$ QDs with and without Eu dopant in GC. b) TEM image of CsPbBr$_3$ QDs without Eu dopant in glass (quantum dot size distribution shown in inset). c) HRTEM image of CsPbBr$_3$ GC. d) TEM image of CsPbBr$_3$:Eu GC (quantum dot size distribution shown in inset). e) Transmittance spectra of CsPbBr$_3$:Eu QDs in glass as a function of Eu concentration (transmittance spectrum of PG measured as a reference). f) Photographs of CsPbBr$_3$ GC, CsPbBr$_3$:1.5% Eu GC (optimal concentration of Eu dopant) and conventional CsPbBr$_3$ nanocrystal film with thickness of about 100 µm (CsPbBr$_3$ nanocrystals and polystyrene are uniformly mixed with the ratio of 1:5). g) Pb 4f XPS spectra of CsPbBr$_3$ GC with and without Eu dopant. h) Eu 3d$_{5/2}$ XPS spectrum of CsPbBr$_3$:1.5% Eu GC. i) Photoluminescence (PL) spectra of CsPbBr$_3$:Eu GC as a function of Eu concentration under 365 nm UV excitation (photograph of CsPbBr$_3$:1.5% Eu GC emission and enlargement of Eu$^{3+}$ emission shown in the inset). PL spectra is measured with integrating sphere to ensure fair comparison.

conventional CsPbBr$_3$ nanocrystal film under the same circumstance.

To access its potential as X-ray scintillator, we conducted the radioluminescence related experiments. Figure 2a gives the RL spectra of CsPbBr$_3$ GC with and without Eu dopant under X-ray excitation. The RL intensity of CsPbBr$_3$ GC at 530 nm enhances significantly with Eu$^{3+}$ doping, consistent with the PL results. Different from the PL spectra, the relatively strong characteristic emissions of Eu$^{3+}$ with peaks at 595, 616, 654, and 701 nm, which can be assigned to $5D_0 \rightarrow 7F_J$ ($J = 1, 2, 3, 4$)$^{[46,47]}$ transitions of Eu$^{3+}$ ions respectively, become more distinct under X-ray excitation. This feature brings an extra advantage that the overlap between RL spectra and absorption is largely reduced leading to decreased self-absorption,$^{[48]}$ while the stokes shift of mere CsPbBr$_3$ perovskite is very small due to its direct bandgap nature. The integrated RL intensity of CsPbBr$_3$:Eu GC as a function of incident X-ray dose rate is shown in Figure 2b revealing a super-linear relation, which is beneficial for obtaining good X-ray image contrast. The discrepancy between PL and RL for CsPbBr$_3$:Eu GC is possibly associated with the different excitation mechanism under UV and X-ray. The incident X-ray photons first interact with the heavy atoms of CsPbBr$_3$ via photoelectric effect and Compton
Characterizations of CsPbBr₃:Eu GC as X-ray scintillator. a) Radioluminescence (RL) spectra of CsPbBr₃ GC, CsPbBr₃:1.5% Eu GC (optimal concentration of Eu dopant) under X-ray excitation with a dose rate of 189 µGy/s at a voltage of 50 kV. The absorption spectra is also included to demonstrate the self-absorption issue. b) RL intensity dependence on X-ray irradiation. c) X-ray attenuation efficiency of CsPbBr₃ GC, CsPbBr₃:1.5% Eu GC, commercial Csl:Tl, and LuAG:Ce. d) Steady-state light yield calculation of our optimized CsPbBr₃:1.5% Eu GC, commercial LuAG:Ce, inset figure is the measurement setup with integrating sphere. e) Afterglows of CsPbBr₃ GC scintillators with and without Eu dopant, compared with widely-used Csl:Tl commercial scintillator. f) Photograph of CsPbBr₃:1.5% Eu GC under X-ray illumination (dose rate of 189 µGy/s, voltage: 50 kV).

scattering. Note that the molar ratio of CsPbBr₃ and Eu₂O₃ (Eu³⁺ source) is ≈17:1.5, thus CsPbBr₃ QDs are the major contributors of X-ray absorption as shown in Figure 2c. Once the X-ray photons are absorbed, they will be first converted to high-energy ejected electrons. These electrons with high kinetic energies move and scatter though their surrounding materials, loosing energy progressively by causing other secondary electrons, and then are transported to the luminescent centers (CsPbBr₃ and Eu³⁺ ions in this case). In our case, we suppose that many of the hot and secondary electrons originated from CsPbBr₃ are captured by the Eu³⁺ luminescent centers since the Eu³⁺ are incorporated into CsPbBr₃ crystal lattice. This conjecture can be confirmed by comparing the RL intensity of reference glass sample with the same amount of 1.5% Eu and CsPbBr₃:1.5% Eu GC, in which the Eu³⁺ RL intensity is more intense than the reference glass sample as shown in Figure S4, Supporting Information. Figure 2d gives the estimation of the X-ray light yield of the perovskite glass-ceramic scintillator under steady state X-ray illumination, the reference sample was LuAG:Ce which had similar light decay constant to perovskite scintillators. It demonstrated the CsPbBr₃:1.5% Eu GC scintillator had an estimated X-ray light yield (steady-state) of 10 100 photons per MeV. The light yield describes the X-ray to photon conversion efficiency, and is considered as an internal quantum efficiency, an analogy to PL quantum yield. Hence its value significantly depends on how the X-ray absorption is determined, which is sort of ambiguous here. Herein we treated the whole perovskite GC as the X-ray absorber although the glass matrix was non-emissive, in such a way that the value of 10 100 photons per MeV was obtained. The ultimate output light intensity depends on not only light yield but also X-ray absorption coefficient and light outcoupling efficiency. Fast light decay of perovskite scintillator is considered as one of the major advantages over the conventional scintillators. One might have the concern of whether this unique feature is still kept with Eu doping. The results in Figure 2e relieve us that the CsPbBr₃:1.5% Eu GC has almost as fast light decay as non-doped CsPbBr₃, both clearly outperforming commercial Csl:Tl. The measurement was conducted with an industrial test kit in Hamamatsu scintillator department, with which the X-ray source can be immediately shut off by instantly shifting the electron beam off the metal target. The time correlated single photon counting measurement in Figure S5, Supporting Information, indicates the decay time of CsPbBr₃:1.5% Eu GC and CsPbBr₃ GC are 6.78 and 2.89 ns respectively, both in the nanosecond regime. Finally, Figure 2f shows the photographs of CsPbBr₃:Eu GC scintillations under X-ray excitation exhibiting strong and uniform emissions. Table S2, Supporting Information, illustrates the figures of merits of perovskite glass ceramic scintillator, together with previously reported perovskite scintillators and some typical commercial scintillators. Despite decades of intensive research of
Figure 3. Spatial resolution evaluation of X-ray imaging based on three scintillators: CsPbBr₃ nanocrystal film, CsPbBr₃ GC, and CsPbBr₃:1.5%Eu GC.

a) The schematic of X-ray imaging system. b) Photograph of chip A. c) X-ray images of chip A using CsPbBr₃ GC scintillators with and without Eu³⁺ doping (dose rate: 47.2 µGy air s⁻¹, voltage: 50 KV). d) Modulation transfer functions (MTF) of X-ray images obtained from CsPbBr₃ GC scintillators with and without Eu³⁺ doping, the conventional CsPbBr₃ nanocrystal film scintillator is set as reference. The image’s spatial resolutions (when MTF value equals 0.2) acquired from CsPbBr₃ GC, CsPbBr₃:Eu GC, and CsPbBr₃ QDs scintillators are 15.0, 4.1, and 1.5 lp mm⁻¹, respectively. e) Bright field and X-ray image of standard X-ray resolution test pattern plate using CsPbBr₃:Eu GC scintillator (dose rate: 189 µGy air s⁻¹, voltage: 50 KV).

Inorganic scintillators, there is not yet a single material that combines all the figures of merits, instead people have developed many scintillators for different applications. The perovskite GC scintillator, though not ideal, has balanced parameters of moderate light yield and fast light decay, which might be useful in some of the real applications, for example, dynamic X-ray imaging or even medical CT that requires high time resolution and high stability.

The above characterizations imply the potential of CsPbBr₃:Eu GC as scintillator for X-ray imaging. Herein, we built a homemade X-ray imaging system in Figure 3a, and successfully acquired X-ray images of chip A (optical photograph shown in Figure 3b) by using CsPbBr₃ GC without and with Eu³⁺ doping. As expected, the image quality using CsPbBr₃:Eu GC is clearly better than the one using non doped CsPbBr₃ GC, due to the enhanced RL intensity and reduced light scattering (Figure 3c). The X-ray image of chip B acquired by conventional CsPbBr₃ nanocrystal film reveals quite weaker image sharpness because of the strong optical crosstalk induced by light scattering (Figure S6, Supporting Information). The modulation transfer functions (MTF)⁵¹ of images obtained from CsPbBr₃ GC, CsPbBr₃:Eu GC, and CsPbBr₃ nanocrystal film have been calculated by the slanted-edge method⁵² to compare the spatial resolution (Figure 3d). X-ray edge images used for MTF calculation are shown in Figure S7, Supporting Information. The spatial resolution is defined to be the spatial frequency (lp mm⁻¹) at MTF = 0.2. As a result, the image resolution with CsPbBr₃:Eu GC scintillator is 15.0 lp mm⁻¹ which is much larger than 4.1 lp mm⁻¹ for CsPbBr₃:GC, while conventional CsPbBr₃ nanocrystal film only gives 1.5 lp mm⁻¹. To further confirm those values, we took images of the standard X-ray resolution test pattern plate (Figure 3e), showing the observation limit was between ≈14 and 16 lp mm⁻¹, consistent with its calculated MTF value. The corresponding X-ray images of standard test pattern plate for referencing CsPbBr₃ GC and CsPbBr₃ nanocrystal film are presented in Figure S8, Supporting Information. It is noted that the referencing CsPbBr₃ nanocrystal film was reported with 9.8 lp mm⁻¹ in the system with scintillator screen very close to the photodiode array,⁵¹ such that the optical crosstalk of scattered scintillation light makes minimum impact. In our case, the camera is placed away from the scintillator screen in order to increase field of view, and we should expect even better resolution if the transparent GC scintillator is attached closely to the CMOS sensor of same pixel size. Even so, this is still the highest spatial resolution for perovskite-based X-ray imaging⁵¹,⁵² to the best of our knowledge, it is even superior to 10 lp mm⁻¹ of the typical commercial CsI(Tl) scintillator based X-ray imaging.¹⁶

The X-ray radiation stability, thermal and environmental stability of CsPbBr₃:Eu GC scintillator has been systematically investigated in terms of RL, X-ray image quality. The RL intensity...
has retained almost 100% of its initial value following continuous X-ray irradiation for 100 h and thermal annealing for another 100 h under 85 °C, clearly outperforming the conventional nanocrystal film shown in Figure 4a. The relative humidity was also recorded every 10 h showing a humid ambient of the testing environment (RH: ≈40% and 80%). The corresponding X-ray images of chip B were acquired at three specific timing during the aging process (0, 100, and 200 h), which gave nearly identical sharpness visually (Figure 4b). X-ray images of chip B obtained at every 10 h during the stability test are presented in Figure S9, Supporting Information. The stability tests demonstrate robustness of our CsPbBr$_3$:Eu GC as scintillator against harsh ambient environment (high humidity, thermal annealing) and longtime X-ray illumination. Figure 4c shows high-quality X-ray image of human finger with obvious biological tissue phase contrast and clear joint details benefited from the low optical crosstalk of the scintillation. Figure 4d,e are the X-ray images of different circuit boards with various electronic components. To access the suitability of CsPbBr$_3$:Eu GC scintillator for the real-time X-ray imaging, a video-rate X-ray imaging of artificial bone bending (Video S2, Supporting Information) was made which exhibited a distinct phase contrast without ghost imaging effect. Moreover, we monitored the incubation process of a fertilized quail egg by acquiring its X-ray images during the incubation process (Figure S10, Supporting Information). The image obtained at the 12th day showed obvious phase contrast and a pair of wings appeared indistinctly. These X-ray imaging demos demonstrate great potential of CsPbBr$_3$:Eu GC scintillator for the medical or industrial X-ray imaging application.

In conclusion, we have successfully developed an Eu doped perovskite glass-ceramic scintillator to eliminate the scintillation scattering and improve operational stability. The incorporation of Eu promotes the crystallization of CsPbBr$_3$ QDs, and makes them distribute uniformly in the matrix. As a result, high resolution X-ray imaging of 15 lp mm$^{-1}$ is realized. In the meanwhile, the operational and environmental stability of CsPbBr$_3$:Eu GC...
scintillators are significantly enhanced. This approach and platform could also be applicable to other perovskite or perovskite derivative scintillators that might possess better figures of merits than CsPbBr$_3$ scintillators.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

W.M. and T.J. contributed equally to this work. The authors acknowledge the support from the National Key Research and Development Program of China (2017YFA0207700), Outstanding Youth Fund of Zhejiang Natural Science Foundation of China (LR18F050001), National Nature Science Foundation of China (61804134), National Nature Science Foundation of China (NSFC) (61965012), Yunnan Ten Thousand Talents Plan Young & Elite Talents Project (YNWR-QNBJ-2018-295, YNWR-QNBJ-2018-325), the Excellent Young Project of Yunnan Province Applied Basic Research Project (2019FI001).

Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

Research data are not shared.

Keywords

perovskites, quantum dots, scintillator, spatial resolution, X-ray imaging

Received: September 30, 2020
Revised: January 31, 2021
Published online:

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