Discovery of millihertz quasi-periodic oscillations in the X-ray binary EXO 0748–676

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ABSTRACT

We report the discovery of millihertz quasi-periodic oscillations (mHz QPOs) from the bursting, high-inclination atoll neutron star low-mass X-ray binary (NS LMXB) EXO 0748–676 with the Rossi X-ray Time Explorer (RXTE). This class of QPO, originally discovered in three NS LMXBs, has been interpreted as a consequence of a special mode of nuclear burning on the NS surface. Using all the RXTE archival observations of the source, we detected significant (>3σ) mHz QPOs in 11 observations. The frequency of the oscillations was between ∼5 and ∼13 mHz. We also found a decrease of the QPO frequency with time in two occasions; in one of these the oscillations disappeared with the onset of an X-ray burst, similar to what was reported in other sources. Our analysis of the X-ray colours revealed that EXO 0748–676 was in a soft spectral state when it exhibited the QPOs. This makes EXO 0748–676 the sixth source with mHz oscillations associated with marginally stable burning, and the second one that shows a systematic frequency drift. Our results suggest that the mechanism that produces the drift might always be present if the mHz QPOs are observed in the so-called intermediate state.

Key words: accretion, accretion discs – stars: individual: EXO 0748–676 – stars: neutron – X-rays: binaries.

1 INTRODUCTION

Low-mass X-ray binaries (LMXBs) are some of the brightest X-ray sources in the Galaxy. These systems consist of a compact object (a neutron star or a black hole) and a `normal' (generally a main-sequence ≲1 M⊙) companion star. These LMXBs are powered by mass accretion, from the donor star to the compact object, through Roche lobe overflow. Because of the conservation of angular momentum, this material does not fall directly on to the compact star but forms an accretion disc around it (Pringle & Rees 1972).

Neutron star (NS) LMXBs are classified into two main groups based on their timing and spectral properties, the Z and atoll sources, the latter being fainter than the former (Hasinger & der Klis 1989). In turn, the atoll sources are divided in three main spectral states: hard (or extreme island), soft (or banana), and intermediate (or island) states. NSs in LMXBs are usually thought to be weakly magnetized (B < 1010 G; e.g. Psaltis & Chakrabarty 1999 and references therein). Many of these systems show X-ray bursts, a sudden and very rapid increase in the X-ray flux due to a thermonuclear runaway, followed by an exponential decay due to the cooling of the NS (see e.g. Strohmayer & Bildsten 2006). Most of the NS LMXBs show X-ray variability in the form of relatively sharp mHz to kHz quasi-periodic oscillations (QPOs), and broad-band noise at <100 Hz (see e.g. review by van der Klis 2006). Typical well-defined classes of QPOs are, among others, the so-called (i) kHz QPOs, which generally appear in pairs in the 300–1300 Hz range (see e.g. van der Klis 2006), (ii) hectohertz (hHz) QPOs, with frequencies in the range of ∼100–300 Hz (e.g. van Straaten, van der Klis & Méndez 2003, Altamirano et al. 2008b), (iii) low-frequency QPOs, which have frequencies in the 0.01–50 Hz range.
The mHz QPOs have properties that set them apart from the other types of QPOs: mHz QPOs occur only within a narrow range of X-ray luminosities \((L_{\text{2-20keV}} \approx (0.5–1.5) \times 10^{37} \ \text{erg s}^{-1})\), their fractional rms amplitude increases towards low energies, and the QPOs disappear at the onset of a thermonuclear (type I) X-ray burst. Revnivtsev et al. (2001) suggested that the mHz QPOs could correspond to a special mode of nuclear burning on the NS surface, although they could not rule out an explanation due to instabilities in the accretion disc. The results of Yu & van der Klis (2002) further supported the thermonuclear interpretation: they found that in the atoll source 4U 1608–52 the frequency of the kHz QPO was anticorrelated with the 2–5 keV X-ray count rate variations during the mHz oscillation. This anticorrelation suggests that the inner edge of the accretion disc increases due to the additional flux produced in each mHz QPO cycle on the NS surface. This is opposite to the positive correlation with those parameters (Heger, Cumming & Woosley 2007). These authors also found that all the bursts that come up immediately after the mHz QPOs had positive convexity, suggesting that both phenomena happen at the NS equator.

Altamirano et al. (2008a) found a systematical decrease of the mHz QPO frequency with time in the NS LMXB 4U 1636–53. Once the frequency dropped below 9 mHz, a bright type I X-ray burst occurred and the oscillations disappeared. This drifting behaviour was only observed when 4U 1636–53 was close to the transition between the hard and soft states. Altamirano et al. (2008a) showed that the mHz QPO frequency can be used to predict the occurrence of a type I burst, strengthening the interpretation that the mHz QPO phenomenon is intimately related to nuclear burning on the NS surface.

Additional work of Lyu, Méndez & Altamirano (2014) and Lyu et al. (2015, 2016) investigated the properties of the mHz QPOs in 4U 1636–53. Using RXTE and XMM–Newton observations, the authors found that the mHz QPO frequency and the temperature of the NS surface were not significantly correlated. This finding is not consistent with theory, which predicts an anticorrelation between those parameters (Heger, Cumming & Woosley 2007). These authors also found that all the bursts that come up immediately after the mHz QPOs had positive convexity, suggesting that both phenomena happen at the NS equator.

From the theoretical side, Heger et al. (2007) suggested that the mHz QPOs are produced by marginally stable nuclear burning of He on the NS surface. The model of Heger et al. (2007) predicts an oscillatory mode of burning, with a period of \(\approx 100 \ \text{s} \) (similar to what is observed), and only happening near the transition between stable and unstable burning (first predicted by Paczynski 1983), i.e. within a very narrow range of mass accretion rate. However, while in this model the oscillations appear at accretion rates close to the Eddington rate, observationally the QPOs were observed at an order of magnitude lower \(L_b\). Heger et al. (2007) explained the discrepancy by proposing that the accreted material covers only \(\approx 10\) per cent of the NS surface, so that the local accretion rate (accretion rate per unit area) is effectively close to that of Eddington.

The model of Heger et al. (2007) cannot explain all the phenomenology observed (e.g. mHz QPOs and X-ray bursts occurring at the same time and the frequency drift). To further investigate this issue, Keek, Cyburt & Heger (2014) studied the possible relation between the mHz QPOs, the composition of the accreted fuel and the nuclear reaction rates. They found that neither the variation in the composition nor the reaction rate are responsible for the appearance of the mHz QPOs at the observed accretion rate. However, Keek, Langer & in’t Zand (2009) investigated the effect of NS rotation on the stability of the burning. They found that the combination of a turbulent chemical mixing of the accreted fuel and a higher heat flux from the crust increase the stability of He burning, such that it could explain the observed accretion rate at which the mHz QPOs take place. Most importantly, Keek et al. (2009) found that the QPO frequency correlates with the heat flux from the crust, so that the frequency drift observed by Altamirano et al. (2008a) could be due to the cooling of the deeper layers of the NS where the burning occurs.

EXO 0748–676 is a transient NS LMXB discovered in outburst during EXOSAT observations in February of 1985 (Parmar et al. 1985). Its X-ray light curve exhibits \(\approx 8.3\)-min duration eclipses, implying an inclination of \(\approx 75^\circ–82^\circ\) (Parmar et al. 1986). The eclipses allowed to estimate an orbital period of \(3.82 \ \text{h}\). The source also shows irregular X-ray dips, mostly in the half orbital cycle prior to the eclipse (Parmar et al. 1986). EXO 0748–676 also presents type I X-ray bursts, a few of which exhibited photospheric radius expansion (PRE). Depending on the assumed fuel composition, the distance estimated using PRE bursts was between \(d = 5.9 \pm 0.9 \) and \(7.4 \pm 0.9 \ \text{kpc}\) (Wolff et al. 2005; Galloway et al. 2008). Homan & van der Klis (2000) detected a kHz QPO during the persistent emission of an observation of this source with RXTE. Galloway et al. (2010) reported a 552-Hz burst oscillation, associated with the spin frequency of the NS in EXO 0748–676.

2 OBSERVATIONS AND DATA ANALYSIS

We used all pointed observations of EXO 0748–676 obtained with the Proportional Counter Array (PCA; for instrument information see Jahoda et al. 2006) aboard the Rossi X-ray Timing Explorer (RXTE). We studied a total of 749 observations taken over nearly 15 yr (March 1996–September 2010). An observation is usually composed of more than an orbit (each one of \(\approx 90 \ \text{min}\) duration), separated by data gaps of at least 1 ks. In a few occasions the conditions were such that RXTE continuously observed the source for up to 17 ks. Thus, our data sets consist of data segments of length \(\sim 0.4–17 \ \text{ks}\). We used 1 s resolution PCA (observed in Event, Good-Xenon, and Single Bit modes) data to create light curves. Under the assumption that the mHz QPOs are stronger at lower energies, we followed Altamirano et al. (2008a) and searched for them in the \(\approx 2–5 \ \text{keV}\) energy range (channels 0–10), applying the Lomb–Scargle periodogram (Lomb 1976; Scargle 1982; Press et al. 1992) to each gap-free segment separately. In those segments in which one or more type I X-ray bursts were detected, we searched for periodicities before, after and in between bursts. When an eclipse was detected, we searched for variability before and after it. We found that the oscillations in the \(\approx 2–5 \ \text{keV}\) energy band are evident.

\(^{1}\)Linares et al. (2010) reported mHz QPOs from the 11 Hz pulsar IGR J17480–2446, but the properties of the QPOs are quite different compared with those of the listed sources. Ferrigno et al. (2017) found mHz QPOs in the accreting millisecond X-ray pulsar IGR J00291+5934, but the QPO origin is unclear.
from the light curves. In all cases, the detections reported in this Letter are above the 3σ level. Following Altamirano et al. (2008a), we calculate the significance level taking into account the number of frequencies searched and assuming white noise as prescribed in the method of Press et al. (1992). We also measured the frequency of the mHz QPO, νQPO, as the frequency at which the power reaches its maximum value in its respective periodogram.

In order to create the colour–colour and hardness-intensity diagrams, we calculated colours using the 16-s time-resolution Standard 2 mode data and followed the exact same method used by Altamirano et al. (2008b), with the addition that for this source we identified and excluded eclipsing episodes as well as type I X-ray bursts. We defined the soft colour as the 3.0–6.0 keV/2.0–3.5 keV count rate ratio, the hard colour as the 9.7–16.0 keV/6.0–9.7 keV count rate ratio, and the intensity as the 2.0–16.0 keV count rate. All values were normalized by the colours of the Crab in the same energy bands.

3 RESULTS

In Fig. 1 we show both the 1.5–12 keV monitoring 15-d average light curve from the All-Sky Monitor (ASM; Levine et al. 1996) and the 2–16 keV light curve obtained from RXTE/PCA pointed observations of EXO 0748–676. The intensity of the source was approximately constant, from MJD ∼ 50000 to MJD ∼ 53000, after which a number of X-ray flares occurred, and EXO 0748–676 transitioned to the quiescent state.

In Fig. 2 we show the colour–colour diagram (CCD) of EXO 0748–676. This plot resembles the CCD of the atoll source 4U 1636–53 (Altamirano et al. 2008a), even if the softest branch is missing (so-called banana branch). We identified the possible extreme island state (EIS) with the hardest colours (which is the most populated part of the diagram), the island state (IS) with intermediate colours, and the lower left banana (LLB) state with the softest colours. This last identification is supported by the detection of kHz QPOs (Homan & van der Klis 2000), which are generally observed in this state (e.g. review by van der Klis 2000). A reliable identification of the states would require a detailed variability study (e.g. van Straaten et al. 2003, Altamirano et al. 2008b) which is beyond the scope of this Letter.

We detected mHz QPOs with frequencies between ∼5.3 and ∼12.8 mHz in 14 data segments (all shorter than 4 ks, distributed in 11 observations) during the flaring period between 2004 March and 2007 December (MJD 53071 and 54439, respectively; see Fig. 1). All detections were found when the source was between ≈1 per cent and ≈3 per cent Edd, assuming a distance of 7.1 kpc and a 3–50 keV $L_{\text{3–50 keV}} = 2.5 \times 10^{38} \text{ erg s}^{-1}$.

The mHz QPOs (filled red squares; see Fig. 2) appear when the source has hard colour $\lesssim 0.9$ and soft colour $\lesssim 1$, that is, probably the LLB. We also found mHz QPOs in two segments that fall in an intermediate state, close to the transition between the banana and the island states. In all cases the 2–5 keV fractional rms amplitude of the QPOs was ∼4 percent, and the QPOs were not significantly detected for energies higher than ∼7 keV, consistent with the results of Revnivtsev et al. (2001).

In Fig. 3 we show a representative example of a background subtracted light curve of two consecutive observations (obsIDs 80040-01-11-00, –01; from now on, observation ‘B’; see Figs 1 and 2). The oscillations can be seen from the beginning of the first observation to ∼10.5 ks until an X-ray burst occurred, after which the QPOs disappeared. The QPOs were not detected after the burst for 11.3 k until the end of that observation. The disappearance of the mHz QPOs after the X-ray burst is consistent with what Revnivtsev et al. (2001) and Altamirano et al. (2008a) found for other sources, and suggests that the physical processes producing the QPOs and the bursts are connected. In the inset of Fig. 3 we show the periodogram of the first data segment. A prominent QPO at ∼8 mHz and a possible harmonic at ∼16 mHz can be seen.

Altamirano et al. (2008a) found a systematic drift in the frequency of the mHz QPOs of some observations of the NS LMXB 4U 1636–53. The length of the segments where we detected mHz QPOs
QPOs exhibited the frequency drifts, consistent with what Altamirano et al. (2008a) found in 4U 1636−53. In later work, Linares et al. (2012) reported that these mHz QPOs appear as the result of a smooth evolution of the thermonuclear bursts properties: as the luminosity of the source increased (decreased), the burst rate increased (decreased), and the fluence of each burst decreased (increased), so that at luminosities higher than $L_{2–50	ext{keV}} \sim 10^{38} \text{erg s}^{-1}$, bursts appeared in the light curve as a QPO with a period of $\sim 3 \text{min}$. These results suggested that the burst rate and burst fluence are a function of the luminosity (accretion rate). This behaviour is inconsistent with our findings, where the mHz QPO and the type I X-ray bursts occur at the same observed luminosity. However, it is currently under debate if the phenomenology observed in IGR J17480−2446 and the six sources listed above are produced by different physical processes, or it is the same physical process under different conditions (see discussion in Linares et al. 2012). For example, different composition of the burning material, different accretion geometry (i.e. accretion mainly on to the NS equator versus spherical accretion on to most of the NS surface), different NS spin frequency and/or magnetic field, and a combination of all these, can affect the burning process and how we detect it.

Ferrigno et al. (2017) detected a QPO at an approximately constant frequency of $\sim 8 \text{mHz}$ in the accreting millisecond X-ray pulsar (AMXP) IGR J00291+5934. The QPO was consistently observed in the background corrected light curve ($\sim 2–5 \text{keV}$) of EXO 0748−676 rebinne to $32 \text{s}$ in which the mHz QPOs are detected (obs ‘B’; label as in Figs 1 and 2). The oscillations are present in three segments, two before and one after the eclipse. About $10.5 \text{ks}$ from the start of the first observation, a type I X-ray burst occurred and the QPOs, apparent before the burst, disappeared. The inset shows the Lomb–Scargle periodogram of the first $\sim 1200 \text{s}$ of the data. A significant QPO is present at a frequency of $8 \text{mHz}$, and a second harmonic component at $\sim 16 \text{mHz}$ appears in the plot.

**Figure 3.** RXTE/PCA light curve ($\sim 2–5 \text{keV}$) of EXO 0748−676 rebinned to $32 \text{s}$ in which the mHz QPOs are detected (obs ‘B’; label as in Figs 1 and 2). The oscillations are present in three segments, two before and one after the eclipse. About $10.5 \text{ks}$ from the start of the first observation, a type I X-ray burst occurred and the QPOs, apparent before the burst, disappeared. The inset shows the Lomb–Scargle periodogram of the first $\sim 1200 \text{s}$ of the data. A significant QPO is present at a frequency of $8 \text{mHz}$, and a second harmonic component at $\sim 16 \text{mHz}$ appears in the plot.

**Figure 4.** Background corrected light curve ($\sim 2–5 \text{keV}$) of the observation ‘A’ (label as in Figs 1 and 2) of EXO 0748−676 taken with RXTE/PCA, using $10\text{s}$ time bins. QPOs are clearly visible through the whole observation. are generally too short to measure frequency drift; however in two cases we were able to combine data to test for drifts. The first case is shown in Fig. 3. The QPO frequency drifts from segment to segment, starting at $\sim 8.2 \text{mHz}$ (first segment), and drifting to $\sim 5.8$ and $\sim 5.4 \text{mHz}$, before and after the eclipse, respectively. Assuming that the QPOs were always present until the burst onset, the frequency decreased at an average rate of $\sim 0.26 \text{mHz ks}^{-1}$. The second case is from obsID 90039-01-03-05 (three orbits; from now on, observation ‘A’; see Figs 1 and 2), where the QPO frequency drifts from about $\sim 13 \text{mHz}$ to less than $6 \text{mHz}$. Fig. 4 shows the light curve and Fig. 5 the dynamical power spectrum of this observation. The QPOs in this observation persisted for about $13 \text{ks}$ (assuming continuity during data gaps), and their frequency drifted at an average rate of $\sim 0.56 \text{mHz ks}^{-1}$.

EXO 0748−676 was close to the intermediate state when the QPOs exhibited the frequency drifts, consistent with what Altamirano et al. (2008a) found in 4U 1636−53. Finally, excluding those cases where we observed a drift, the frequency of the mHz QPOs were between $\sim 6.7$ and $\sim 8.7 \text{mHz}$. Our results, therefore, show that EXO 0748−676 is the sixth source that displays this kind of QPO and the second one where a frequency drift is observed.

**Figure 5.** Dynamical power spectrum corresponding to the RXTE observation ‘A’ (label as in Figs 1 and 2) of EXO 0748−676. The plot, which was made using a $700 \text{s}$ sliding window with steps of $50 \text{s}$, shows the evolution of a frequency component associated with the mHz QPOs, decreasing from $\approx 13 \text{mHz}$ down to $\approx 6 \text{mHz}$.

**4 DISCUSSION**

We report the discovery of mHz QPOs in the NS LMXB EXO 0748−676. The QPO frequency ranged between $\sim 5$ and $\sim 13 \text{mHz}$, its amplitude was $\sim 4$ per cent at low energy ($\sim 2–5 \text{keV}$), and in at least two cases the QPO frequency decreased in time. All the detections occurred when the source energy spectrum was relatively soft, as suggested by the CCD. In one case we can confirm that the QPO disappeared after the onset of a thermonuclear X-ray burst (in none of the cases we saw the QPO disappear without the onset of a burst).

We report the discovery of mHz QPOs in the NS LMXB EXO 0748−676. The QPO frequency ranged between $\sim 5$ and $\sim 13 \text{mHz}$, its amplitude was $\sim 4$ per cent at low energy ($\sim 2–5 \text{keV}$), and in at least two cases the QPO frequency decreased in time. All the detections occurred when the source energy spectrum was relatively soft, as suggested by the CCD. In one case we can confirm that the QPO disappeared after the onset of a thermonuclear X-ray burst (in none of the cases we saw the QPO disappear without the onset of a burst). All these observational properties are consistent with those observed in the accreting NS LMXBs 4U 1636−53, 4U 1608−52, and Aql X−1 (Revnivtsev et al. 2001; Altamirano et al. 2008a), 4U 1323−619 (Strohmayer & Smith 2011), and GS 1826−238 (Strohmayer et al. 2018, although in this source the amplitude of the mHz QPO increased with energy), suggesting that we are observing the same phenomenology in six different accreting systems. These mHz QPOs have been associated with an oscillatory mode of He nuclear burning on the NS surface, generally referred as marginally stable nuclear burning (e.g. Heger et al. 2007; Kekelidze et al. 2009).

Linares et al. (2010) reported the detection of mHz QPOs with a frequency of $\sim 4.5 \text{mHz}$ during the 2010 outburst of the NS transient IGR J17480−2446. In later work, Linares et al. (2012) reported that these mHz QPOs appear as the result of a smooth evolution of the thermonuclear bursts properties: as the luminosity of the source increased (decreased), the burst rate increased (decreased), and the fluence of each burst decreased (increased), so that at luminosities higher than $L_{2–50\text{keV}} \sim 10^{38} \text{erg s}^{-1}$, bursts appeared in the light curve as a QPO with a period of $\sim 3 \text{min}$. These results suggested that the burst rate and burst fluence are a function of the luminosity (accretion rate). This behaviour is inconsistent with our findings, where the mHz QPO and the type I X-ray bursts occur at the same observed luminosity. However, it is currently under debate if the phenomenology observed in IGR J17480−2446 and the six sources listed above are produced by different physical processes, or it is the same physical process under different conditions (see discussion in Linares et al. 2012). For example, different composition of the burning material, different accretion geometry (i.e. accretion mainly on to the NS equator versus spherical accretion on to most of the NS surface), different NS spin frequency and/or magnetic field, and a combination of all these, can affect the burning process and how we detect it.
4.1 Frequency drifts and source state

Altamirano et al. (2008a) found that the mHz QPOs in 4U 1636–53 would show clear frequency drifts if the system was observed close to the transition between the island and banana states. If the QPOs were detected in the banana state, then the QPO frequency did not drift. Our results on EXO 0748–676 are consistent with those of Altamirano et al. (2008a): in the only two instances in which we observe frequency drifts (A and B, see Fig. 2), EXO 0748–676 was in an intermediate state.

The QPOs in observation ‘A’ showed a faster frequency drift than we observed in observation ‘B’ (~0.56 mHz ks\(^{-1}\) versus ~0.26 mHz ks\(^{-1}\), respectively). In terms of spectral states, observation ‘A’ is at a harder state than ‘B’. Under the assumption that in the banana state we would not see any frequency drift (as reported by Altamirano et al. 2008a, for 4U 1636–53), our results appear to suggest that the rate at which the QPO frequency decreases is related to the spectral state, where the rate is maximum at the hardest colours, and decreases as the source evolves to softer states, where we do not expect to see any drift at all. This is not a scenario that can be tested with our RXTE data of EXO 0748–676; however, it could be tested with other sources (e.g. in 4U 1636–53, where Altamirano et al. 2008a reported 22 instances of mHz QPOs with frequency drifts).

Finally, our results show that the frequency drift is not a unique characteristic of the marginally stable burning in 4U 1636–53, but that can be seen in other sources. This opens a window of opportunity to potentially better understand this type of burning. For example, comparison of the mHz QPOs characteristics between sources could shed light into the dependence of the QPO frequency drift, and the change in heat flux from the NS crust (Keek et al. 2009).

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