The stellar populations of galaxies hold vital clues to their formation histories. In this paper we present results based on modeling stacked spectra of early-type galaxies drawn from the Sloan Digital Sky Survey as a function of velocity dispersion, \(\sigma\), from 90 km s\(^{-1}\) to 300 km s\(^{-1}\). The spectra are of extremely high quality, with typical signal-to-noise ratio of 1000 Å\(^{-1}\), and a wavelength coverage of 4000 Å − 8800 Å. Our population synthesis model includes variation in 16 elements from C to Ba, a two-component star formation history, the shift in effective temperature, \(\Delta T_{\text{eff}}\), of the stars with respect to a solar metallicity isochrone, and the stellar initial mass function, among other parameters. In our approach we fit the full optical spectra rather than a select number of spectral indices and are able to, for the first time, measure the abundances of the elements V, Cr, Mn, Co, and Ni from the integrated light of distant galaxies. Our main results are as follows: (1) light-weighted stellar ages range from 6–12 Gyr from low to high \(\sigma\); (2) [Fe/H] varies by less than 0.1 dex across the entire sample; (3) Mg closely tracks O, and both increase from \(\approx 0.0\) at low \(\sigma\) to \(\approx 0.25\) at high \(\sigma\); Si and Ti show a shallower rise with \(\sigma\), and Ca tracks Fe rather than O; (4) the iron peak elements V, Cr, Mn, and Ni track Fe, while Co tracks O, suggesting that Co forms primarily in massive stars; (5) C and N track O over the full sample and [C/Fe] and [N/Fe] exceed 0.2 at high \(\sigma\); and (6) the variation in \(\Delta T_{\text{eff}}\) with total metallicity closely follows theoretical predictions based on stellar evolution theory. This last result is significant because it implies that we are robustly solving not only for the detailed abundance patterns but also the detailed temperature distributions (i.e., isochrones) of the stars in these galaxies. A variety of tests reveal that the systematic uncertainties in our measurements are probably 0.05 dex or less. Our derived [Mg/Fe] and [O/Fe] abundance ratios are 0.05–0.1 dex lower than most previous determinations. Under the conventional interpretation that the variation in these ratios is due to star formation timescale variations, our results suggest longer star formation timescales for massive early-type galaxies than previous studies. Detailed chemical evolution models are necessary in order to translate the abundance ratio distributions of these galaxies into constraints on their formation histories. Alternatively, these data may provide useful constraints on the nucleosynthetic pathways for elements whose production is not well understood.

Key words: galaxies: abundances – galaxies: stellar content

Online-only material: color figures
The sensitivity of [α/Fe] to all of these processes significantly complicates the interpretation of this ratio. Nonetheless, it has become standard practice to assume that the [α/Fe] ratio is telling us something about the timescale of star formation. By comparing to simple closed-box chemical evolution models, Thomas et al. (2005) concluded that the [α/Fe] ratios of the most massive galaxies in their sample, with σ ≈ 300 km s^{-1}, implied star formation timescales of only ~200 Myr. These short inferred star formation timescales, if correct, would have profound implications for our understanding of the formation of the most massive galaxies in the universe.

The analysis of spectral indices is beginning to give way to the modeling of the full optical–NIR spectra of stellar systems. This approach has great potential, and is now frequently used to measure star formation histories and metallicities non-parametrically (Heavens et al. 2000; Cid Fernandes et al. 2005; Ocvirk et al. 2006; Tojeiro et al. 2009). Recently, such full-spectrum fitting models have been extended to include variation in the elemental abundance patterns (Walcher et al. 2009; Conroy & van Dokkum 2012a), paving the way to robustly measure many parameters from integrated light spectra.

This paper is part of an ongoing series focused on modeling very high quality optical–NIR spectra of early-type galaxies with full-spectrum fitting stellar population synthesis (SPS) techniques. In Conroy & van Dokkum (2012a, CvD12), we presented the model, and in Conroy & van Dokkum (2012b) we used that model to measure the low-mass IMF from the spectra of 34 nearby early-type galaxies and the nuclear bulge of M31 (based on data presented in van Dokkum & Conroy 2011). In Conroy et al. (2013) we presented results on the neutron-capture elements Sr and Ba based on our local galaxy sample and on the stacked spectra that are the focus of this work. In the present paper we measure the ages, detailed abundance patterns, and effective temperatures of the stars in early-type galaxies drawn from the Sloan Digital Sky Survey (SDSS). In a subsequent paper we will discuss the low-mass IMFs inferred for these galaxies.

The rest of this paper is organized as follows. Section 2 provides an overview of the SPS model and our fitting technique, Section 3 describes the data, and Section 4 presents a test of the model by fitting to spectra of metal-rich star clusters. Our main results are presented in Section 5, and in Section 6 we compare our derived properties to results based on other modeling techniques. A discussion and summary of our results is provided in Sections 7 and 8.

2. MODEL AND FITTING TECHNIQUE

The SPS model used herein was developed in Conroy & van Dokkum (2012a, hereafter CvD12), to which we refer the interested reader for details. The model employs standard SPS techniques, including combining libraries of isochrones and stellar spectra. We use isochrones from three separate groups, all of which are solar metallicity with solar-scaled abundance patterns.

Empirical spectra form the base of the model. We make use of two separate libraries, the MILES library, which covers the wavelength range 0.35 μm < λ < 0.74 μm (Sánchez-Blázquez et al. 2006), and the IRTF library of cool stars, which covers the wavelength range 0.81 μm < λ < 2.4 μm (Cushing et al. 2005; Rayner et al. 2009). The empirical stars are of approximately solar metallicity and have solar abundance patterns. We have computed a large grid of synthetic stellar atmospheres and spectra in order to model the relative change in the spectrum of a star due to a change in the abundance of a single element. We chose 20 positions along a 13 Gyr isochrone to compute synthetic spectra for each abundance pattern. The model atmospheres and spectra were computed with the ATLAS12 model atmosphere and spectrum synthesis package (Kurucz 1970, 1993), ported to Linux by Sbordone et al. (2004). The line list was provided by R. Kurucz.4 including linelists for TiO and H2O, among other molecules. The synthetic stellar spectra were grafted onto the empirical stellar spectra in order to create models with arbitrary abundance patterns. We emphasize that the synthetic spectra are only used differentially. The model allows for arbitrary variation in the IMF and spans ages from 3–13.5 Gyr. Figure 1 shows the model spectrum for a 13 Gyr, solar metallicity population, smoothed to a velocity dispersion of σ = 350 km s^{-1}.

There have been several updates to the CvD12 model. Most importantly, the synthetic atmospheres and spectra have been recomputed with an updated line list kindly provided by R. Kurucz. This line list includes many more iron lines and updated wavelengths and hyperfine splitting of many lines. The model now contains variation in many more elements, described below. We have also supplemented the MILES empirical library with stacked SDSS M dwarf spectra from Bochanski et al. (2007). This turns out to be important because the MILES library contains only a few M dwarf spectra, and none cooler than M6V. The SDSS M dwarf template spectra cover the entire M dwarf spectral sequence. The final update since CvD12 is the inclusion of emission lines in the fitting (rather than masking regions of potential emission line contamination, as done previously). Currently we include 13 emission lines, including individual components of [O iii] λ5007, [Ne iii] λ3868, Hα, Hγ, Hβ, [O iii] λ5007, [N i] λ5201, [N ii] λ6547, 6583, Hα, and [S ii] λ6716, 6731. We include separate parameters for the strengths of each of these lines (i.e., we do not enforce particular line ratios, even among doublets with known, constant line ratios). The intrinsic line shape is assumed to be a gaussian with a single width for all lines. This width is also included in the fit.

The effect of varying elemental abundances on the model spectra are shown in Figures 2–4. These figures show the response of a 13 Gyr, solar metallicity model to an increase in the abundance of a single element. We emphasize that the model atmospheres and spectra are always treated self-consistently (i.e., the elemental concentration is increased both in the computation of the model atmosphere and of the model spectra). In many cases the effect on the spectrum from a change in an element is straightforward, e.g., increasing [Mg/H] results in stronger Mg i lines, but other effects are more subtle. For example, because Mg is a major electron donor, increasing [Mg/H] affects the ionization equilibrium and therefore affects many species of other elements. For example, it is for this reason that [Mg/H] affects the strength of the Ca ii lines at ~0.86 μm.

The effect of varying age from 7 to 13 Gyr is shown in Figure 5. A variety of temperature-sensitive features are evident, including the hydrogen balmer lines, NaD, and the Mg i lines. Most importantly, comparison of Figure 5 to Figures 2–4 shows that when one considers information in the full spectrum, age and metallicity effects are readily separable.

We follow Conroy & van Dokkum (2012b) in fitting the model to data. In its present form the model contains 40 free parameters, including the redshift and velocity dispersion, a two-part

\footnote{kurucz.harvard.edu/}
Figure 1. Model spectrum for an age of 13 Gyr and solar metallicity. The spectrum has been smoothed with a velocity dispersion of $\sigma = 350$ km s$^{-1}$, equal to the smoothing applied to the early-type galaxy data analyzed in this paper. Strong features are labeled. Also included is the location of the true stellar continuum, which is the spectrum that would be observed in the absence of all line opacity. In this figure the model spectrum is computed entirely from synthetic stellar spectra, whereas for the main analysis the synthetic spectra are only used differentially.

Figure 2. Relative response of the blue spectral region to an increase in the abundance of nine different elements. Abundances have been increased by 0.3 dex for all elements except for C, which is increased by 0.15 dex. The changes are with respect to a solar metallicity model with a MW IMF and an age of 13 Gyr. The models have been smoothed to a velocity dispersion of $\sigma = 150$ km s$^{-1}$. Notice the different y-axis range in each panel. These are noiseless spectra, and so every feature visible in the figure is real. It is evident from both this figure and Figures 3 and 4 that essentially every spectral region is influenced by one or more elements.

power-law IMF, two population ages (the age of the dominant population and the mass fraction of a 3 Gyr population), four nuisance parameters, 13 emission line strengths, the velocity broadening of the emission lines, and the abundances of C, N, Na, Mg, Si, Ca, Ti, V, Cr, Mn, Fe, Co, Ni, Sr, and Ba, and O, Ne, S (the latter three are varied in lock-step). These parameters are fit to the data via a Markov Chain Monte Carlo (MCMC) fitting technique. The data and models are split into four wavelength intervals (described below) and, within each interval the spectra are normalized by a high-order polynomial (with degree $n$ where
\[ n \equiv (\lambda_{\text{max}} - \lambda_{\text{min}}) / 100 \, \text{Å}; \text{ see Conroy \& van Dokkum 2012b, for details).} \]

For the purposes of the present article, the emission line strengths are treated as additional nuisance parameters, and in many cases their strengths are constrained to be very low, which is not surprising given the sample definition described in the next section. Unless specified otherwise, ages quoted in this paper are light-weighted stellar ages computed from the two age components.

One of the nuisance parameters included in the model is a shift in effective temperature, \( \Delta T_{\text{eff}} \), applied to all stars used in the synthesis. As described in CvD12, this parameter is meant to capture changes in the isochrones due to changes in abundance patterns of the model stars. The motivation for this shift is shown in Figure 6. In the top panel we show isochrones obtained from the Dartmouth Stellar Evolution Database (DSEP; Dotter et al. 2008) for a range in \([\text{Fe}/\text{H}]\) and \([\alpha/\text{Fe}]\). In the bottom panel we have shifted each isochrone by a constant amount in \( T_{\text{eff}} \) until a match was achieved with the solar metallicity case. A constant shift does an excellent job of accounting for modest variation in abundance patterns from the solar metallicity model (see also Salaris et al. 1993). We caution that simple shifts in \( T_{\text{eff}} \) work less well at low metallicity (VandenBerg et al. 2012). Moreover, it is interesting that the isochrone with \([\text{Fe}/\text{H}] = -0.1 \) and \([\alpha/\text{Fe}] = 0.2 \) is almost identical to the solar metallicity isochrone. Thus, over the metallicity range of interest in this work, a simple shift in \( T_{\text{eff}} \) is probably sufficient to account

\begin{figure}
\centering
\includegraphics[width=\textwidth]{figure3.png}
\caption{Same as Figure 3, now showing the red spectral region. Abundances have been increased by 0.3 dex for all elements except for C, which is increased by 0.15 dex.}
\end{figure}

\begin{figure}
\centering
\includegraphics[width=\textwidth]{figure4.png}
\caption{Same as Figure 2, now showing the elements encompassing the iron peak (V, Cr, Mn, Fe, Co, Ni). The change in the model spectra have been computed for a 0.3 dex increase in the abundance of each element.}
\end{figure}
for elemental abundance effects on the isochrones. Because the abundances of many elements are varied simultaneously in the model, we use the free parameter $\Delta T_{\text{eff}}$ to encompass their combined effect, rather than attempting to compute self-consistent isochrones from first principles. However, we caution that a constant shift in $T_{\text{eff}}$ is unlikely to capture the full variation in isochrone morphology with elemental abundance. For example, a change in $[^{\alpha}/Fe]$ at fixed $[Fe/H]$ results in a slightly cooler main sequence turn-off point and giant branch but a slightly warmer lower main sequence (Dotter et al. 2007; VandenBerg et al. 2012).

3. DATA

3.1. SDSS Data

The main analysis presented here results from fitting our updated CvD stellar population model to very high-signal-to-noise ratio (S/N) stacks of passive galaxies, binned by stellar velocity dispersion. These galaxies are selected from the SDSS (York et al. 2000) Main Galaxy Survey (Strauss et al. 2002) Data Release 7 (Abazajian et al. 2009), within a narrow redshift interval ($0.025 < z < 0.06$). Following the methodology of Peek & Graves (2010), we select passive galaxies by requiring that they have no detected emission in $\text{H}\alpha$ nor in $\text{[O ii]}$ λ3727.

The individual spectra have moderate S/N, typically $\sim 20 \text{Å}^{-1}$, so we stack the spectra of hundreds of galaxies to achieve the very high S/N needed for detailed abundance analysis. Passive galaxy star formation histories vary strongly as a function of their stellar velocity dispersion, but also at fixed $\sigma$ depending on how far they scatter off the Fundamental Plane (FP; Graves et al. 2010; Springob et al. 2012). We therefore have chosen to stack galaxies in bins of $\sigma$, in order to study the main dimension of stellar population variation, but restrict our sample to galaxies that lie on the FP (i.e., in the central FP slice defined in Graves et al. 2010). The two-dimensional space of galaxy star formation histories will be explored in a future paper.

We stack the spectra in seven bins of velocity dispersion, with mean values of $\sigma = 88, 112, 138, 167, 203, 246$, and 300 km s$^{-1}$. The effective resolution of the data therefore varies from $R \approx 400$–1200 owing to intrinsic Doppler broadening. The SDSS spectra are obtained through 3$''$ diameter fibers, which sample the inner $0.8 R_e$ for the smallest $\sigma$ bin and the inner $0.4 R_e$ for the largest $\sigma$ bin. Each individual spectrum was continuum-normalized and convolved to an effective dispersion of 350 km s$^{-1}$ before stacking, and each spectrum contributed equally to the stack. Problematic pixels (e.g., those under bright sky lines) were masked before stacking. The resulting stacked spectra have S/N at 5000 Å that ranges from $\approx 500 \text{Å}^{-1}$ to $\approx 1800 \text{Å}^{-1}$. Four wavelength intervals were defined for the SDSS stacks, 0.40 μm–0.48 μm, 0.48 μm–0.58 μm, 0.58 μm–0.64 μm, and 0.80 μm–0.88 μm for fitting models to data.

In Figure 7 we show the stacked spectra in three $\sigma$ bins. The fluxes have been continuum-normalized in order to focus attention on the differences in the line strengths between different stacks. Notice that the differences are relatively subtle, with features strengths varying by only a few percent from the lowest to the highest dispersion bins.

3.2. Star Cluster Data

In this paper we will also analyze integrated light spectra of metal-rich star clusters in order to test the model. We focus on two $\sim Z_{\odot}$ clusters, M67 and NGC 6528. The former is an intermediate-age open cluster with solar-scaled abundance
ratios, while the latter is an old bulge globular cluster with enhanced abundance ratios. For M67 we adopt an age of 4 Gyr derived from its color–magnitude diagram (CMD; VandenBerg et al. 2007; Magic et al. 2010), [Fe/H] = 0.0, and solar-scaled abundance ratios with an uncertainty of 0.1 dex (Tautvaišienė et al. 2000; Shetrone & Sandquist 2000). For NGC 6528 we adopt an age of 11 Gyr, consistent with recent CMD-based determinations (Feltzing & Johnson 2002; Momany et al. 2003). The elemental abundance data are derived from high-resolution spectra of individual red giant branch (RGB) stars, as reported in Carretta et al. (2001), Zoccali et al. (2004), and Origlia et al. (2005).

The integrated light spectrum for M67 was constructed in Schiavon et al. (2004) by combining spectra of individual member stars with a Salpeter (1955) IMF. M67 contains a large population of blue stragglers, but these were omitted in the integrated spectrum. The spectrum covers the range 3640 Å to 5400 Å, at a resolution of 2.7 Å. The S/N of the integrated spectrum was not provide by Schiavon et al., so we adopt a nominal S/N of 100 per pixel.

For NGC 6528 we use the integrated light spectrum from the library of globular cluster spectra presented in Schiavon et al. (2005). These spectra were obtained via drift scan observations, so the spectra reflect the true integrated light of the clusters. The spectrum covers the wavelength range 3360–6430 Å at a resolution of ∼3.1 Å. The S/N ranges from ≈50 per pixel at 4000 Å to 150 per pixel at 5000 Å. In this paper we analyze the spectrum extracted within the central 7′′. Several spectral regions had to be masked owing to sub-optimal sky subtraction and bad CCD columns. These regions include 4155 Å–4165 Å, 4535 Å–4565 Å, 4850 Å–4870 Å, and 5035 Å–5065 Å.

For both spectra we use two wavelength intervals when fitting the model to these data: 4000 Å–4600 Å, and 4600 Å–5350 Å.

4. TESTING THE MODEL AGAINST METAL-RICH STAR CLUSTERS

Before discussing the main results of this paper, in this section we present the results of fitting our models to spectra of the metal-rich star clusters M67 and NGC 6528. Star clusters are of course simpler to model than galaxies, owing to the fact that they are at least approximately coeval and monometallic. They therefore can serve as useful tests of any SPS model.

In Figure 8 we compare the best-fit ages, [Fe/H], and several α elements for these two clusters. Turning first to M67, we derive an age of 3.5 ± 0.3 Gyr, in excellent agreement with the CMD-based age. We also find excellent agreement for Fe, Mg, Ca, and Ti. The abundances of O and Si are underestimated by 1σ–2σ. Because the effects of O and Si on the spectrum are somewhat more subtle than the other α elements, integrated light spectra that include the redder spectral coverage for M67 are needed before a more definitive conclusion regarding these elements can be reached. Though not shown in Figure 8, we also find [C/Fe] = −0.06 and [N/Fe] = −0.02 for M67, which agrees within 0.1 dex with literature determinations of C and N from main sequence turn-off stars in M67 (Shetrone & Sandquist 2000).

The model also provides an excellent fit to the age and [Fe/H] abundance of NGC 6528. The α elements also broadly agree with the literature data, although perhaps the most striking impression from Figure 8 is the large scatter between the literature estimates for the abundance ratios. The most dramatic difference is in the estimated [Ca/Fe] abundance, where Origlia et al. (2005) find [Ca/Fe] = +0.4 while Zoccali et al. (2004) quote [Ca/Fe] = −0.40; a difference of 0.8 dex! We have plotted the Origlia et al. (2005) results in grey because these authors model high-resolution NIR spectra while Carretta et al. (2001) and Zoccali et al. (2004) model high-resolution blue spectra. We also draw attention to the fact that these two authors find relatively modest enhancements in O and Mg, less than 0.2 dex. Thus, while NGC 6528 is α-enhanced, it is not tremendously so.

For NGC 6528 we do not compare results for C nor N because these elements can be synthesized in the envelopes of RGB stars, so abundances of these elements derived from such stars do not reflect the original (i.e., main sequence) abundances. Our model recovers the latter, whereas literature estimates of C and N for this cluster are based on RGB stars. Owing to the large discrepancy in reported literature abundances, even for important species such as Fe, O, Mg, Si, Ca, and Ti, we made no effort to compare to other elements as their derived abundances from high-resolution stellar spectroscopy appear to be even more uncertain.

Our model successfully recovers the ages, [Fe/H], and α element abundances of metal-rich star clusters. Further progress on testing SPS models must await (1) higher quality integrated light spectra of star clusters covering a wider wavelength range and at higher S/N than is presently available; and (2) a thorough evaluation of the results from high-resolution stellar spectroscopy. Until a consensus emerges on the abundance patterns of the stars in metal-rich star clusters, such systems cannot be used to rigorously test abundance ratio predictions of SPS models.

5. RESULTS

We now turn to our main results—the analysis of high quality stacked spectra of early-type galaxies from the SDSS. In this section we present and discuss the overall quality of the fits, the detailed abundance patterns, and the derived effective temperatures of the stars. Table 1 presents the main derived parameters from modeling the stacked spectra. Results for the abundance ratios of neutron-capture elements Sr and Ba were presented in Conroy et al. (2013), and the results for the variation
Figure 8. Best-fit parameters for the star clusters M67 (top panel) and NGC 6528 (bottom panel). Results from our full-spectrum fitting of the integrated light (open blue circles) are compared to literature values based on high-resolution spectra of individual member stars (red lines and symbols). Fiducial ages for M67 and NGC 6528 are taken to be 4 and 11 Gyr, respectively, based on isochrone fitting to their CMDs. Abundance patterns for NGC 6528 are from Carretta et al. (2001, C01), Zoccali et al. (2004, Z04), and Origlia et al. (2005, O05). For M67 we adopt a fiducial abundance pattern of $[\text{Fe}/\text{H}] = -0.0 \pm 0.1$ (roughly $1\sigma$ errors Tautvaišienė et al. 2000; Shetrone & Sandquist 2000). Overall the agreement between our model predictions and the literature values is good, although the large variation in the literature determination of many of the elemental abundances for NGC 6528 prohibits a more definitive test of our model.

(A color version of this figure is available in the online journal.)

| log $\sigma$ ($\text{km s}^{-1}$) | log $M_\odot$ ($M_\odot$) | $\text{rms}$ | age (Gyr) | log $f_\gamma$ | [Fe/H] | C | N | O | Na | Mg | Si | Ca | Ti | V | Cr | Mn | Co | Ni |
|-----------------|-----------------|---------|----------|-------------|--------|---|---|---|----|----|----|----|----|---|---|---|---|
| 1.94            | 9.63            | 0.22    | 6.2      | -0.7        | -0.07  | 0.05| 0.02| 0.03| -0.16| 0.05| 0.02| 0.03| 0.05| 0.03| 0.02| 0.03| 0.03| -0.00|
| 2.05            | 9.75            | 0.22    | 6.5      | -0.9        | -0.05  | 0.07| 0.08| 0.06| -0.09| 0.08| 0.00| 0.03| 0.07| 0.01| 0.00| 0.03| 0.06| -0.00|
| 2.14            | 9.80            | 0.21    | 6.8      | -2.2        | -0.04  | 0.09| 0.13| 0.10| 0.02| 0.09| 0.03| 0.03| 0.08| 0.03| -0.01| 0.03| 0.10| 0.01|
| 2.23            | 10.08           | 0.22    | 7.0      | -4.1        | -0.03  | 0.13| 0.18| 0.17| 0.12| 0.12| 0.05| 0.03| 0.09| 0.03| -0.01| 0.03| 0.16| 0.02|
| 2.30            | 10.55           | 0.26    | 7.6      | -4.3        | -0.01  | 0.16| 0.21| 0.20| 0.22| 0.15| 0.09| 0.03| 0.11| 0.03| -0.03| 0.02| 0.20| 0.03|
| 2.39            | 10.70           | 0.29    | 11.0     | -4.3        | -0.02  | 0.19| 0.26| 0.25| 0.33| 0.20| 0.13| 0.04| 0.12| -0.01| -0.02| 0.05| 0.27| 0.01|
| 2.47            | 11.07           | 0.34    | 11.8     | -4.4        | 0.00   | 0.21| 0.27| 0.28| 0.43| 0.22| 0.16| 0.02| 0.12| -0.02| -0.03| 0.05| 0.26| 0.07|

Notes. Stellar masses are averages within the $\sigma$ bins and they assume a Kroupa (2001) IMF. The rms deviation between the model and data is computed over the full wavelength range used in the fit. Ages are light-weighted and the fraction of mass in the 3 Gyr component, $f_\gamma$, is included in the table. Abundances of elements are quoted as $[\text{X}/\text{Fe}]$. Formal statistical errors are less than 1% and so are omitted from the table. Systematic errors dominate the error budget, and are estimated to be 0.05 dex or less. See Section 5.4 for further discussion.

5.1. Overall Quality of the Fits and Uniqueness of the Derived Parameters

We begin by discussing the overall quality of the model fits. In Figures 9 and 10 we show the data and best-fit model for the lowest and highest $\sigma$ bins. Within each panel the fluxes have been continuum-normalized by a high-order polynomial. These figures demonstrate the overall very high quality of the fits. In the lowest dispersion bin the rms between data and model is only 0.22%, and $\chi^2/\text{dof} = 1.1$, indicating that our model is able to fully capture the observed variation in the spectrum to the precision allowed by the data. At the highest dispersion bin, where the abundance ratios deviate most strongly from the solar-scaled values, the rms is 0.34%, with $\chi^2/\text{dof} = 4.7$. It is not surprising that the fit is formally worse for the highest dispersion bin, both because the S/N is higher, which imposes more stringent demands on the model, and because the strong

in the IMF with galaxy $\sigma$ will be presented in future work. We simply note here that the IMF trends with $\sigma$ are similar to those presented in Conroy & van Dokkum (2012b). We include in this table estimated stellar masses based on the mean $r$-band luminosities of the galaxies in each bin combined with the derived stellar $M/L$ values assuming a Kroupa (2001) IMF.

Table 1

Results from Modeling Stacked SDSS Early-type Galaxy Spectra

| log $\sigma$ ($\text{km s}^{-1}$) | log $M_\odot$ ($M_\odot$) | $\text{rms}$ | age (Gyr) | log $f_\gamma$ | [Fe/H] | C | N | O | Na | Mg | Si | Ca | Ti | V | Cr | Mn | Co | Ni |
|-----------------|-----------------|---------|----------|-------------|--------|---|---|---|----|----|----|----|----|---|---|---|---|
| 1.94            | 9.63            | 0.22    | 6.2      | -0.7        | -0.07  | 0.05| 0.02| 0.03| -0.16| 0.05| 0.02| 0.03| 0.05| 0.03| 0.02| 0.03| 0.03| -0.00|
| 2.05            | 9.75            | 0.22    | 6.5      | -0.9        | -0.05  | 0.07| 0.08| 0.06| -0.09| 0.08| 0.00| 0.03| 0.07| 0.01| 0.00| 0.03| 0.06| -0.00|
| 2.14            | 9.80            | 0.21    | 6.8      | -2.2        | -0.04  | 0.09| 0.13| 0.10| 0.02| 0.09| 0.03| 0.03| 0.08| 0.03| -0.01| 0.03| 0.10| 0.01|
| 2.23            | 10.08           | 0.22    | 7.0      | -4.1        | -0.03  | 0.13| 0.18| 0.17| 0.12| 0.12| 0.05| 0.03| 0.09| 0.03| -0.01| 0.03| 0.16| 0.02|
| 2.30            | 10.55           | 0.26    | 7.6      | -4.3        | -0.01  | 0.16| 0.21| 0.20| 0.22| 0.15| 0.09| 0.03| 0.11| 0.03| -0.03| 0.02| 0.20| 0.03|
| 2.39            | 10.70           | 0.29    | 11.0     | -4.3        | -0.02  | 0.19| 0.26| 0.25| 0.33| 0.20| 0.13| 0.04| 0.12| -0.01| -0.02| 0.05| 0.27| 0.01|
| 2.47            | 11.07           | 0.34    | 11.8     | -4.4        | 0.00   | 0.21| 0.27| 0.28| 0.43| 0.22| 0.16| 0.02| 0.12| -0.02| -0.03| 0.05| 0.26| 0.07|

Notes. Stellar masses are averages within the $\sigma$ bins and they assume a Kroupa (2001) IMF. The rms deviation between the model and data is computed over the full wavelength range used in the fit. Ages are light-weighted and the fraction of mass in the 3 Gyr component, $f_\gamma$, is included in the table. Abundances of elements are quoted as $[\text{X}/\text{Fe}]$. Formal statistical errors are less than 1% and so are omitted from the table. Systematic errors dominate the error budget, and are estimated to be 0.05 dex or less. See Section 5.4 for further discussion.
Figure 9. Comparison between the SDSS early-type galaxy stacked spectrum in the lowest velocity dispersion bin (black lines) and the best-fit model (red lines). Within each wavelength range, the top panel shows the continuum-normalized fluxes and the lower panel shows the ratio between data and model. The grey band denotes the 1σ error limits of the data. Overall the quality of the fit is excellent.

(A color version of this figure is available in the online journal.)

Figure 10. Same as Figure 9, now showing results for the highest velocity dispersion bin. As in Figure 9 the overall quality of the fit is excellent.

(A color version of this figure is available in the online journal.)
deviations from solar-scaled abundance ratios means that the synthetic stellar spectra play an important role in the modeling. The synthetic spectra carry significant uncertainties (e.g., due to the treatment of the model atmospheres and incomplete line lists), so their greater weight in the model will tend to result in larger model uncertainties.

In Figure 11 we show the residuals now for all seven bins. The most striking feature of the residuals is their similarity across the sample. For example, the residuals around the Ca ii triplet at 8500 Å–8700 Å are relatively significant, but the magnitude of the residuals does not vary strongly with galaxy velocity dispersion, indicating that our model is not capturing the totality of variation in the observed spectra. In addition, deviations from solar-scaled abundance ratios means that the synthetic stellar spectra play an important role in the modeling. The synthetic spectra carry significant uncertainties (e.g., due to the treatment of the model atmospheres and incomplete line lists), so their greater weight in the model will tend to result in larger model uncertainties.

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demonstrates that the response functions of these elements are similar in several respects. This similarity is driven in part by molecular equilibrium involving CO and TiO. But from Figure 12 it is evident that our model is capable of separately constraining C, O, and Ti very well (notice the small range of the x- and y-axes).

(A color version of this figure is available in the online journal.)

Figure 12. Covariance between [O/Fe], [C/Fe] and [Ti/Fe] for the SDSS stacked spectrum with log $\sigma = 2.39$ km s$^{-1}$. Contours show 68% and 95% confidence limits in the posterior distribution. The response function due to O variation is similar (though opposite in sign) to C variation at $\lambda < 6000$ Å (see Figure 2), which is mostly a consequence of molecular equilibrium involving CO. Similarly, in the red, the response functions of O and Ti resemble one another, here because of TiO. This figure demonstrates that, in the context of our model, we are able to separately constrain C, O, and Ti very well (notice the small range of the x and y-axes).

5.2. Abundance Patterns

We now turn to the elemental abundance ratios derived from our sample of stacked SDSS early-type galaxy spectra.

In Figure 13 we show the relations between velocity dispersion and C and N (left panel), and the $\alpha$ elements O, Mg, Si, Ca, and Ti (right panel). In most cases the formal statistical errors are smaller than the symbol sizes. The abundance of both C and N increase from [C/Fe] $\approx$ [N/Fe] $\approx 0.0$ at low $\sigma$ to 0.2–0.3 at high $\sigma$. This is in qualitative agreement with previous results (Kelson et al. 2006; Graves et al. 2007; Schiavon 2007; Smith et al. 2009; Johansson et al. 2012).

The derived abundance ratio patterns of the $\alpha$ elements show several interesting features. First, all ratios are close to the solar-scaled values at low $\sigma$. Second, the trends with $\sigma$ are stronger for the lighter elements O and Mg compared to the heavier elements Si, Ca, and Ti. Mg appears to track O closely, with a slight preference for [Mg/O] $< 0.0$ at high $\sigma$. Ca tracks Fe closely over the full sample, in agreement with previous work (Saglia et al. 2002; Cenarro et al. 2003; Thomas et al. 2003a; Graves et al. 2007; Schiavon 2007; Smith et al. 2009; Johansson et al. 2012). This suggests that a substantial amount of Ca is formed in Type Ia SNe, as expected from theoretical yields (Nomoto et al. 1984). The enhancement of Si and Ti at high $\sigma$ also broadly agrees with earlier work (Milone et al. 2000; Johansson et al. 2012; Worthey et al. 2013).

In Figure 14 we show the results for [Fe/H] in addition to the other iron peak elements V, Cr, Mn, Co, and Ni, all of which are measured here for the first time. The [Fe/H] abundance increases modestly over the sample by only 0.07 dex. This remarkable uniformity of the Fe abundance, in light of the substantial variation in the abundance of the light and $\alpha$ elements, poses a challenge to chemical evolution models of early-type galaxies. The abundances of the iron peak elements V, Cr, Mn, and Ni closely track Fe over the full sample, implying that these elements form in the same nucleosynthetic sites. In stark contrast is Co, which increases in relative abundance to Fe as $\sigma$ increases, reaching [Co/Fe] $\approx 0.27$ at the highest dispersions.

This result is explored further in Figure 15, where we show the abundance ratios of several elements with respect to O as a function of $\sigma$. Clearly both N and Co track O remarkably closely. C also tracks O more closely than Fe, but [C/O] is not quite as constant as [N/O] and [Co/O], suggesting that C and N may not be forming in exactly the same environments.

(A color version of this figure is available in the online journal.)

Figure 13. Abundance ratios of C, N, and the $\alpha$ elements as a function of early-type galaxy velocity dispersion. In most cases the statistical errors are smaller than the symbol sizes. C, N, O, and Mg all seem to track each other fairly closely, while the heavier $\alpha$ elements Si, Ca, and Ti show weaker trends with $\sigma$. Ca clearly tracks Fe instead of O. Systematic errors are probably $< 0.05$ dex (see Section 5.4).

Figure 14. The relations between velocity dispersion and the elements O, Mg, Si, Ca, and Ti as a function of $\sigma$.
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Figure 14. Abundance ratios of the iron peak elements as a function of early-type galaxy velocity dispersion. In most cases the statistical errors are smaller than the symbol sizes. All iron peak elements track Fe except for Co, which seems to track the light α elements O and Mg. Systematic errors are probably <0.05 dex (see Section 5.4).

(A color version of this figure is available in the online journal.)

Figure 15. Abundance ratios between C, N, Co, and O as a function of early-type galaxy velocity dispersion. All of these elements seem to track each other very closely, especially N, O, and Co. Systematic errors are probably <0.05 dex (see Section 5.4).

(A color version of this figure is available in the online journal.)

Although not shown in a figure, the [Na/Fe] abundance also varies considerably across the sample, from −0.80 to 0.43, qualitatively tracking the light α elements (see Table 1). In detail however the Na abundance variation is considerably stronger than the other elements. The only two Na lines available to us are the NaI features at 5895 Å and 8190 Å, respectively. The former is well-known to be affected by the presence of an interstellar medium (ISM). However, the sample was selected to have no detectable emission in Hα nor in [O II], suggesting that contamination by ISM absorption should be small. Moreover, one might have expected ISM contamination to be greater at lower σ, because low-mass galaxies are younger, and yet the derived [Na/Fe] at low σ is lower than either [O/Fe] or [Mg/Fe] (the presence of an ISM would cause stronger absorption at 5895 Å and thus larger derived [Na/Fe] values). Of course, it is not necessary that [Na/Fe] track the other elements in detail, as its nucleosynthetic origins are quite complex, including contributions from both massive stars and intermediate-mass asymptotic giant branch (AGB) stars. Indeed, the chemical evolution model of Arrigoni et al. (2010) finds a variation in [Na/Fe] in broad agreement with what we find herein, ranging from −0.2 at low mass to +0.3 at high mass.

5.3. Effective Temperatures

As discussed in Section 2, one of the free parameters in our model is the shift in $T_{\text{eff}}$ from a fiducial solar metallicity isochrone. This is an important parameter to include because, as we saw in the previous section, the metallicities and abundance patterns vary by factors of several across our sample. A change in metallicity and/or abundance pattern will induce a change in $T_{\text{eff}}$ of the stars (Dotter et al. 2007). It would be difficult to include this variation self-consistently owing to the large grid of isochrones that would be required. The parameter $\Delta T_{\text{eff}}$ is meant to capture all of these effects into a single number.

In Figure 16 we show the resulting relation between $\Delta T_{\text{eff}}$ and the sum of [O/Fe] and [Fe/H], which is a proxy for [Z/H] (e.g., Trager et al. 2000). The relation derived for our sample of SDSS early-type galaxies is compared to theoretical expectations from the Dartmouth isochrones (Dotter et al. 2008). For the latter, all of the α elements track each other. The early-type galaxy sample increases in σ from low to high metallicity. The theoretical predictions are based on the models shown in Figure 6. The star symbol corresponds to the model with [Fe/H] = −0.1 and [$\alpha$/Fe] = +0.2 and highlights the fact that $\Delta T_{\text{eff}}$ is not a simple function of [Z/H].

The excellent agreement in the slope of the $\Delta T_{\text{eff}}$ versus [O/Fe]+[Fe/H] relation between the theoretical prediction and the data suggests that we are reliably measuring the distribution of effective temperatures for the stars in these galaxies (i.e., we are solving for the detailed shape of the isochrones). The flattening of the observed trend at high σ is interesting and will be explored in detail in future work. The offset between the two
measure the distribution of other by 40 K. The excellent agreement in the varying simultaneously across the sample, and so it is impossible work is ongoing.

relations is $\approx 40$ K, which is well within both the systematic uncertainties in the color-temperature relations used to assign temperatures to empirical stars and the theoretical uncertainties in modeling RGB temperatures. In order to provide a more direct comparison between the observations and the theoretical predictions we will need to create isochrones with the exact same chemical composition as we derive for the early-type galaxies (e.g., variation in C and N in addition to the $\alpha$ elements). This work is ongoing.

We remind the reader that a multitude of parameters are varying simultaneously across the sample, and so it is impossible at present to cleanly separate correlation from causation. The correlation seen in Figure 16 could be driven by a third variable, and the possibility remains that the $\Delta T_{\text{eff}}$ parameter is soaking up an unidentified model deficiency that varies with $\sigma$ and hence $[Z/H]$.

5.4. Exploration of Systematic Uncertainties

In this section we explore variations in the fiducial model and in our primary set of stacked spectra in order to quantify systematic uncertainties in our main results. In Figure 16 we plot the difference in either abundance ratio or log(age) between our main results and a modification to our standard analysis. In the upper left panel we compare our standard analysis to a simplified model with fewer free parameters. In particular, in this model the IMF is fixed to the Kroupa (2001) form, only a single age component is included (rather than the standard two components), the iron peak elements are forced to track Fe, and all the nuisance parameters are turned off (including $\Delta T_{\text{eff}}$). This model results in a larger rms residual between data and model compared to our fiducial model (i.e., it is a poorer fit), which should not be surprising given that there are fewer free parameters. Nonetheless, the resulting abundance ratios agree well with the standard model, in particular Mg, C, Ca, Si, Ti, and Fe agree to within 0.1 dex and in most cases to within 0.05 dex. The largest outliers are the age at low $\sigma$ and the $[O/Fe]$ and $[N/Fe]$ abundances. The difference in $[O/Fe]$ is particularly strong. We speculate that the larger range in $[O/Fe]$ in the simple model is due to the fact that the parameter $\Delta T_{\text{eff}}$ is not included. In essence, $[O/Fe]$ is forced to vary more in order to compensate for the lack of variation in $\Delta T_{\text{eff}}$.

In the upper right panel we consider a model where the mean wavelength-dependent residuals averaged over the sample are added back into the model. In other words, in this case the model is multiplied by the mean residuals shown in Figure 11 before the model is fit to the data. This will obviously result in a much lower rms difference between the best-fit model and data, and the question is whether or not this changes the derived parameters. As demonstrated in Figure 17, the resulting change in abundance patterns is very small, typically $< 0.03$ dex. This implies that the small mis-match between the data and our model does not bias the resulting best-fit parameters.

In the lower left panel we consider the change in parameters when only the blue spectral region ($\lambda < 5800$ Å) is included in the fit. Here again most parameters change very little, by $< 0.05$ dex. Significant outliers include $[Na/Fe]$ and $[O/Fe]$. The change in $[Na/Fe]$ is not surprising because the two features most sensitive to Na are the Na i features at 5895 Å and 8190 Å. With these features masked out the constraint on the Na abundance is coming exclusively from the effect of Na on the electron pressure in the stellar atmospheres, which indirectly affects the ionization states of other elements. For $[O/Fe]$ the deviation is also significant, exceeding 0.1 dex in several bins. The origin of this offset is not clear, although we note that the best-fit $\Delta T_{\text{eff}}$ parameter is lower in the case of fitting only the blue spectra. With cooler stars the metal-lines are stronger, perhaps...
alleviating the need for higher [O/Fe] values. This highlights the need for spectra that extend into the red, where the $\Delta$Fe$_{\text{eff}}$ parameter has a large effect and hence can be more robustly measured. Moreover, at $\lambda < 5800$ Å the constraint on [O/Fe] is coming primarily through the effect of O on molecular dissociation equilibrium in the stellar atmospheres, while at $\lambda > 5800$ Å the effect of O on the spectrum is primarily to increase the strength of the TiO lines (see Figures 2 and 3). The red spectrum therefore offers at the very least a complementary, and arguably a stronger constraint on the [O/Fe] abundance.

In the lower right panel we show the results of fitting to the unsmoothed stacked spectra. This is primarily a test of the model in the low $\sigma$ bins, where the unsmoothed spectra have a resolution $\approx 4 \times$ higher than the smoothed spectra. Here again essentially all parameters agree with our standard results to better than 0.03 dex. The largest outlier at low $\sigma$ is [O/Fe], but even in this case the disagreement is $< 0.1$ dex.

We have performed an additional test (not shown in the figure), where we masked wavelengths bluer than 4200 Å. This was done primarily to test the [Co/Fe] result, since there are several strong Co I lines with hyperfine splitting components at $\lambda < 4200$ Å and many weaker blends at longer wavelengths (see Figure 2). The differences in the derived [Co/Fe] ratios between this and our fiducial model is $< 0.03$ dex for all $\sigma$ bins except for the lowest bin, where the difference is 0.046 dex. The result for Co seems insensitive to the particular wavelength range used in the fit.

In summary, the tests performed in this section suggest that the systematic uncertainties in our derived parameters is probably $< 0.05$ dex.

6. COMPARISON TO PREVIOUS WORK

In this section we provide a comparison to other techniques for deriving stellar population parameters of early-type galaxies. In Figure 18 we show abundance ratios and light-weighted ages from several groups, including the results presented here, from the EZ$_{\text{Ages}}$ code (Graves & Schiavon 2008), and from Thomas et al. (2010, T10), Johansson et al. (2012, J12), and Worthey et al. (2013, W13).

The EZ$_{\text{Ages}}$ code fits 7 Lick indices with the SPS models of Schiavon (2007) in order to derive the age, [Fe/H], [C/Fe], [N/Fe], [Mg/Fe], and [Ca/Fe]. The code was run on the exact same stacked spectra analyzed herein, so that any differences can be attributed entirely to differences in modeling techniques. In general the agreement is very good. The [C/Fe] and [N/Fe] trends are stronger with EZ$_{\text{Ages}}$ compared to the results presented here. Perhaps the most significant differences are the overall higher [Mg/Fe] abundances in EZ$_{\text{Ages}}$ (by $\approx 0.07$ dex), and the lower [Fe/H] abundances (by $\approx 0.07$ dex).

The results from T10 are based on the SPS models of Thomas et al. (2003b) and Thomas et al. (2004). T10 analyzed a morphologically selected sample of early-type galaxies from SDSS. They used 25 Lick indices in order to derive three parameters: [Z/H], age, and [\alpha/Fe]. We compare their derived [\alpha/Fe] values to both [Mg/Fe] and [O/Fe]. We find overall satisfactory agreement between their [\alpha/Fe] values and our [Mg/Fe] and [O/Fe] values, and also between their ages and ours.

J12 present results from the exact same morphologically selected sample of early-type galaxies analyzed in T10, but with the updated SPS models from Thomas et al. (2011). These models allow for variation in the elements C, N, O, Mg, Ca, and Ti, in addition to metallicity and age. They fit their models to the data using 18 Lick indices. The differences between J12 and our own work is the largest of any of the techniques shown in Figure 18 (except for [O/Fe] from W13), although the differences rarely exceed 0.1 dex. It is difficult to know if the differences are due to modeling techniques or sample selection, especially at low $\sigma$ where there is likely to be significant scatter between morphological (“early-type”) and emission line (“quiescent”) selections.

Very recently, W13 analyzed stacked SDSS early-type galaxy spectra using their SPS models. The stacks are not identical to those analyzed herein, but are very similar and a direct comparison between W13 and our results is possible. From Figure 18 it is clear that the ages, [Mg/Fe], [N/Fe], and [Ca/Fe] trends are in good agreement between the techniques, except for the age in the highest $\sigma$ bin in W13. The [C/Fe] abundances are also in reasonable agreement. In contrast, the [O/Fe] trend derived by W13 is much steeper than not only our results but also those of J12 (and T10 if one treats their [\alpha/Fe] as [O/Fe]). In addition, their inferred [Fe/H] abundances range from $-0.05$ at low $\sigma$ to $-0.7$ at high $\sigma$. The results for [Fe/H] are so different from the other results in Figure 18 that they were omitted from the figure for clarity. W13 fixed [Ti/Fe] $= 0.0$ throughout their analysis. These authors also presented results for Si, finding [Si/Fe] is approximately 0.0 at low $\sigma$ and increases to 0.5 at high $\sigma$. This is also in contrast to our results, in which [Si/Fe] increases from 0.0 to 0.16 from low to high $\sigma$. At present we do not know the origins of these discrepancies.

One of the more notable features of Figure 18 is the [Mg/Fe] versus $\sigma$ panel, where there is a $\sim 0.1$ dex variation in the normalization of the relation. This variation is significant in the context of using [Mg/Fe] to infer a star formation timescale (e.g., Thomas et al. 2005). Trager et al. (2008) noted that their models produced [Mg/Fe] values lower than previous Lick index-based models, and they speculated that this may be due to the different treatment of the response functions. In nearly all previous work on Lick index modeling, three theoretical stellar spectra were combined to estimate the response of the integrated light spectrum to abundance variations (Tripicco & Bell 1995; Korn et al. 2005). Trager et al., and also Worthey et al. (2013), instead used many stellar spectra covering the isochrone in order to estimate response functions, in a manner similar to what is done in the CvD models. In addition, previous Lick index-based models calculated response functions for an isochrone with an age of 5 Gyr, whereas more recent models use older isochrones (or a range of isochrones appropriate for the ages of the models).

We have estimated the effects of these assumptions by modifying our model to employ only 3 stars (as opposed to the usual 20) to estimate response functions, and also to use stars along a 5 Gyr isochrone rather than a 13 Gyr isochrone. The number of stars used to compute response functions has a significant effect. We have created response functions for the Mg $b$ Lick index due to a factor of two change in [\alpha/Fe]. A model using 20 stars has a response function in the Mg $b$ equivalent width (EW) that is 30% greater than a model that uses only 3 stars. The effect of the age is rather modest, with a change in EW of only 5%, with 30% greater than a model that uses only 3 stars. The effect of simultaneous abundance variations (Tripicco 1989) was similar to what is done in the CvD models. In addition, previous Lick index-based models calculated response functions for an isochrone with an age of 5 Gyr, whereas more recent models use older isochrones (or a range of isochrones appropriate for the ages of the models).

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Figure 18. Comparison between different techniques for estimating ages and elemental abundances. Our results are compared to techniques that rely on Lick index fitting, including Graves & Schiavon (2008, EZ_Ages), Thomas et al. (2010, T10), Johansson et al. (2012, J12), and Worthey et al. (2013, W13). The EZ_Ages results are based on modeling the exact same stacked spectra as used herein, the W13 results are based on similar but not identical stacks, while the T10 and J12 results are linear fits to results obtained for ~4000 SDSS early-type galaxies with light-weighted ages greater than 2.5 Gyr. Note that T10 derive a single \([\alpha/Fe]\) ratio, which we plot in both the \([Mg/Fe]\) and \([O/Fe]\) panels, as it is not entirely clear which element their \([\alpha/Fe]\) parameter is most sensitive to. Qualitatively the derived trends agree well between various methods/groups, although quantitative differences are clearly evident.

(A color version of this figure is available in the online journal.)

7. DISCUSSION

The main results from this paper are summarized in Figure 19, where we show the abundance patterns of 13 elements as a function of early-type galaxy velocity dispersion. This is the most comprehensive analysis of the detailed abundance patterns of the stars within distant galaxies to-date. We emphasize that the spectra sample approximately the inner 0.5 Re, and thus the abundances derived should be representative of the inner regions of the galaxies. In this section we discuss several implications of these results.

In summary, there is broad agreement between models for those elements that impart strong, localized changes in low-resolution data (especially C, N, Mg, and Ca; see Figure 2). In such cases models based on spectral indices seem to perform as well as our model based on full-spectrum fitting. The power of fitting the detailed spectrum lies in measuring parameters that either impart a more subtle change in the spectrum (as in the case of Si and Mn), or affect large swaths of the spectrum (as in the case of O, Ti, and most of the iron peak elements).
and the strong variation in [Mg/Fe], [C/Fe], and [N/Fe]. It should not be surprising that there is good agreement between techniques for the elements C, N, Mg, and Ca because these elements impart strong, localized changes in low resolution spectra (see Figure 2). Quantitatively there are differences in the slope and normalization for some of these relations, though the differences rarely exceed 0.1 dex in amplitude. Based on our preliminary investigations, it appears that at least some of the differences can be traced back to the number of theoretical spectra used to estimate response functions to elemental abundance variations (see also Trager et al. 2008). The real power of fitting the detailed spectrum is in the measurement of parameters that either impart a small change in the spectrum (e.g., Mn, Sr, Ba), or that affect large regions of the spectrum (e.g., O, Ti, V, Cr, Mn, Fe, Co, Ni). For an example of the former, see Conroy et al. (2013).

The standard interpretation of [α/Fe] is that this ratio is sensitive to the timescale of star formation, with higher values corresponding to shorter timescales (e.g., Tinsley 1979; Thomas et al. 1999). As mentioned in the Introduction, this ratio is also sensitive to the IMF, the nucleosynthetic yields, the delay time distribution of Type Ia SNe, and the preferential loss of metals via winds. Some chemical evolution models are able to reproduce the observed variation in not only Mg and O but also Si and Ca based on the assumption that the main variable driving these trends is indeed the star formation timescale (Pipino et al. 2009), while others also require variation in other parameters such as the IMF (Arrigoni et al. 2010). In the context of any chemical evolution model, the [Mg/Fe] ratio will be correlated with the star formation timescale, and so the lower ratios that we derive herein have important implications for the inferred formation timescales of the stars within these massive galaxies. As an example of the effect, Thomas et al. (2005) use the equation \([\alpha/Fe] \approx (1/5) - (1/6)\log \Delta t\), calibrated from the simple chemical evolution model of Thomas et al. (1999), in order to convert their measured \([\alpha/Fe]\) ratios into constraints on the formation histories of early-type galaxies. In our analysis we find \([Mg/Fe] = 0.22\) in the most massive galaxies, which, when employing the equation above, corresponds to a timescale of 0.8 Gyr. In contrast, the EZ AGES model and T10 find \([Mg/Fe] \approx 0.27\) at the highest dispersions, corresponding to a timescale of 0.4 Gyr. Based on the models of Thomas et al. (2011), J12 find \([Mg/Fe] \approx 0.33\) at the highest dispersions, corresponding to a remarkably short timescale of 0.2 Gyr. It is interesting in this context that the \([O/Fe]\) ratios derived by J12 are in much better agreement with our results at the highest dispersions, and so if \([O/Fe]\) were used as a tracer of \(\alpha\), rather than Mg, then the inferred timescales would be more similar between different groups at high \(\sigma\). However, this of course implies that \([O/Mg] \neq 0\), and in this case it is not clear that a single parameter “[\(\alpha/Fe]\)” can be used to simply (or reliably) translate derived abundance ratios into star formation timescales. Clearly sophisticated chemical evolution models are need to provide further insight.

One of the more interesting trends evident in Figure 19 (see also Figure 15) is the fact that the iron peak element Co closely tracks O rather than Fe. This is the first measurement of Co in the integrated light spectra of early-type galaxies, and it provides an important constraint on the nucleosynthetic origin of Co. Taken at face value, \([Co/Fe] \approx [O/Fe]\) implies that Co forms primarily in massive stars, in apparent contrast to the other iron peak elements. In this context it is worth recalling that Co is also peculiar among the iron peak elements in the metal-poor Galactic halo (McWilliam 1997). This may be a coincidence at some level, as the abundance pattern of the very metal poor halo is probably due exclusively to very metal poor Type II SNe, in contrast to the processes operating to shape the abundance patterns in massive early-type galaxies. Core-collapse SNe nucleosynthesis models predict that Co is synthesized by complete Si-burning in the deepest layers, while Cr and Mn are produced in the outer incomplete Si-burning layers (Woosley & Weaver 1995; Nakamura et al. 1999). There is thus a mechanism by which to separate the production of Co from the other iron peak elements. Perhaps more significantly,

**Figure 19.** Summary of the abundance trends derived herein from stacked spectra of SDSS early-type galaxies. The sample spans a range in velocity dispersion from \(\sigma = 90 \text{ km s}^{-1}\) to \(\sigma = 300 \text{ km s}^{-1}\). The results for Na are omitted from this summary diagram, though we note that [Na/Fe] varies from -0.16 to +0.43 from low to high \(\sigma\).

(A color version of this figure is available in the online journal.)
the Type Ia SNe yields from Nomoto et al. (1984) show a deficit of Co production compared to other iron peak elements. Again, detailed chemical evolution models will be required to sort out these details.

We reaffirm previous results that the light $\alpha$ elements (e.g., O and Mg) show greater variation than the heavier $\alpha$ elements (e.g., Ca and Ti). We also find in this work that the variation in Si is intermediate between the lighter and heavier elements. At least some of these trends can be understood in the context of Type Ia SNe yields. The Nomoto et al. (1984) models produce a large quantity of Ca (comparable to Cr, Mn, Fe, and Ni), and much more Si than either O or Mg. The modest variation in Ti is somewhat more difficult to understand in this context, because the Nomoto et al. models produce very little Ti (comparable to O and Mg). Whether or not the variation in [Ti/Fe] with galaxy dispersion can be reproduced with standard nucleosynthesis yield tables remains to be seen.

Finally we turn to C and N. As found in previous work, both [C/Fe] and [N/Fe] increase with increasing dispersion (Schiavon 2007; Graves et al. 2007; Smith et al. 2009; Johansson et al. 2012). We find that both of these elements track O and Mg quite closely (in detail C seems to track Mg while N tracks O, but we regard this level of similarity as fortuitous). Johansson et al. (2012) also find that C closely tracks O, with [C/O] ~ 0.0, but they find sub-solar [N/O] values. The fact that C and N are tracking the light $\alpha$ elements is probably telling us either that C and N form mostly in massive stars, or that there is a lower limit to the star formation timescale such that both massive stars and moderately massive AGB stars (which can produce lots of C and N) are always contributing to the enrichment of the ISM (see also Schiavon 2007; Johansson et al. 2012, who come to similar conclusions). The latter interpretation is consistent with the star formation timescales derived using our [Mg/Fe] ratios. Much effort is still needed on the theoretical side to understand the nucleosynthetic origins of these elements in order to use them as cosmological clocks.

8. SUMMARY

In this paper we have used the CVD population synthesis model to fit very high quality stacked spectra of SDSS early-type galaxies spanning a range in velocity dispersion from 90 km s$^{-1}$ to 300 km s$^{-1}$. The novel feature of our approach is fitting the full continuum-normalized optical–NIR spectra (4000 Å – 8800 Å), rather than a select number of indices. The model includes variation in age, the abundances of 16 elements, the effective temperature distribution of the stars, the stellar IMF, among other parameters. The resulting quality of the fits are very high, with rms residuals ranging from 0.2%–0.3% across the sample. A variety of tests reveal that the systematic uncertainties in our measurements are probably 0.05 dex or less. We now summarize our main results.

1. Across the sample from low to high $\sigma$, ages increase from 6 to 12 Gyr, [Fe/H] increases from $-0.07$ to 0.00, and the abundance ratios [O/Fe], [Na/Fe], [Mg/Fe], [Si/Fe], [Ti/Fe], and [Co/Fe] increase from approximately zero to 0.1–0.4, with O, Na, Mg reaching the highest values. In contrast, [Ca/Fe], [V/Fe], [Cr/Fe], [Mn/Fe], and [Ni/Fe] remain nearly constant as $\sigma$ increases.

2. N and Co track O to better than 0.03 dex over the full sample, and C tracks O to better than 0.1 dex. Co therefore probably forms predominantly in massive stars. For N and C the situation is more complex because these elements probably form in both massive stars and AGB stars. The constant values of [C/O] and [N/O] therefore either require C and N to form mostly in massive stars or the star formation timescale to never be shorter than several hundred Myr, the lifetime of a $\sim3 M_\odot$ star.

3. Our model allows us to measure the shift in the effective temperature of the stars with respect to a solar metallicity isochrone. This shift inversely correlates with the total metallicity, with a slope that is in excellent agreement with theoretical expectations. We are therefore able to measure not only the mean ages and abundance patterns but also the distribution of effective temperatures of the stars in early-type galaxies.

4. Comparison of our results to Lick index-based techniques reveals agreement in several quantities, such as age, [N/Fe], and [Ca/Fe], and less agreement in others including [C/Fe], [Fe/H], and [Mg/Fe]. Some of the disagreement may be due to different samples of early-type galaxies, but the different modeling techniques also seems to play an important role. Perhaps most relevant from the standpoint of broader implications, we find lower [Mg/Fe] values at fixed $\sigma$ (by 0.05–0.1 dex) than some previous work (especially those which employ the Korn et al. 2005, response tables), which we believe is due to our more accurate modeling of spectral variations due to abundance changes. This implies longer formation timescales for massive early-type galaxies than previously reported.

These results will set the foundation for future studies aimed at constraining the evolutionary histories of galaxies from their detailed abundance patterns. With such a large number of elements precisely measured over a wide range in galaxy mass, it should be possible to place novel constraints not only on the formation histories of these galaxies but also on the nucleosynthetic origins of elements that are difficult to model from first principles.

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