Mechanotransductive cascade of Myo-II-dependent mesoderm and endoderm invaginations in embryo gastrulation

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Animal development consists of a cascade of tissue differentiation and shape change. Associated mechanical signals regulate tissue differentiation. Here we demonstrate that endogenous mechanical cues also trigger biochemical pathways, generating the active morphogenetic movements shaping animal development through a mechanotransductive cascade of Myo-II medio-apical stabilization. To mimic physiological tissue deformation with a cell scale resolution, liposomes containing magnetic nanoparticles are injected into embryonic epithelia and submitted to time-variable forces generated by a linear array of micrometric soft magnets. Periodic magnetically induced deformations quantitatively phenocopy the soft mechanical endogenous snail-dependent apex pulsations, rescue the medio-apical accumulation of Rok, Myo-II and subsequent mesoderm invagination lacking in snail mutants, in a Fog-dependent mechanotransductive process. Mesoderm invagination then activates Myo-II apical accumulation, in a similar Fog-dependent mechanotransductive process, which in turn initiates endoderm invagination. This reveals the existence of a highly dynamic self-inductive cascade of mesoderm and endoderm invaginations, regulated by mechano-induced medio-apical stabilization of Myo-II.
Both biochemical patterning and biomechanical morphogenesis contribute to shaping living organisms during embryonic development. The role of biochemical and genetic information in the development of a functional three-dimensional tissue morphology is beginning to be well understood in different developmental contexts. The reverse role of mechanical cues developed by biomechanical morphogenetic movements, in the activation of developmental biochemical pathways involved in the triggering of cell differentiation, is emerging as a key feature of embryonic development regulation. Mechanical cues generated by developmental strain deformation were found to activate major developmental differentiation events, including early anterior endoderm and mesoderm specification, antero-posterior axis formation, tooth formation, and joint bone differentiation.

Mechanical signalling has also been suggested to trigger the active morphogenetic movements, including mesoderm invagination, folding in legs and wings during organogenesis, primitive streak formation in chick embryos, and cell-cell junction re-enforcement in animal development, as well as meristem morphogenesis, and increased growth rates in plant development. Strains developed by one morphogenetic movement were also suggested to mechanically influence another, for example, with posterior endoderm invagination promoting vertex resolution during germ-band extension. However, the tools used so far to test mechanical triggering of active morphogenetic movements in the context of morphogenetic deficiencies were based on the use of either soft but non-physiological direct mechanical deformation (such as indentation or pipette aspiration) or of genetically induced physiological deformation, which is not directly mechanical (Myosin-II activation in neighbouring cells). This opened the possibility of a response to non-physiological mechanical strains or to uncontrolled indirect biochemical pathways independent of mechanical strains, for instance of a Myo-II-triggered re-enforcement of cell-cell junctions potentially leading to Myo-II recruitment in neighbouring cells in a purely biochemical process, respectively.

By developing an innovative micro-magnet-based tool capable of producing dynamic physiological direct mechanical strain in vivo with cell scale resolution, here we show mechanical triggering of active morphogenetic movements by the mechanical cues induced by preceding morphogenetic movements. Specifically, apical constriction of mesodermal cells, and the resulting mesoderm invagination, which corresponds to the first morphogenetic movement of early Drosophila embryogenesis, are both demonstrated to be mechanically activated by soft apical pulses of the mesoderm.

In the Drosophila embryo, mesoderm invagination represents the first morphogenetic movement of embryogenesis. This movement is triggered by medio-apical stabilization of Myo-II, leading to the acto-myosin apical constriction of mesoderm apaxes that generates tissue bending and folding. Apical constriction follows two distinct phases. The first phase is characterized by snail (sna)-dependent uncoordinated unstable constricting pulsations, associated with unstable medio-apical coalescence of Myo-II spots. The second phase, which leads to invagination, is characterized by a twist (twi)-dependent coordinated stable constriction, associated with stable and progressive coalescence of medio-apical Myo-II spots (Fig. 1a). In sna mutants, both the uncoordinated and coordinated phases are absent, as is mesoderm invagination, thus revealing the interplay between sna and twi pathways in the regulation of mesoderm invagination. Mid-apical stabilization of Myo-II, collective apical constriction and mesoderm invagination can be mechanically rescued in sna mutants by indenting the mesoderm by a few microns. This stimulates a mechanosensitive pathway, Folded gastrulation (Fog), that is downstream of twist. Such a response to a non-physiological mechanical perturbation suggested that the mechanotransductive mechanical induction of medio-apical stabilization of Myo-II by the physiological mechanical strains resulting from sna-dependent pulsations could regulate the interplay between sna and twi pathways, which in turn leads to mesoderm morphogenesis.

We show that the medio-apical stabilization of Myo-II in early Drosophila embryo mesoderm is mechanotransductively triggered in response to strain developed by sna-dependent pre-existing individual cell apex pulsations. This is demonstrated by mimicking the physiological apical pulsations absent in sna mutants, resulting in the rescue of the Rho-dependent medio-apical stabilization of Rho-associated protein kinase (Rok) and Myo-II, and then mesoderm invagination. Pulsations are induced using time varying, cell scale spatially localized mechanical forces generated within sna mutant embryos, by exploiting magnetic interactions between magnetic nanoparticles (MNPs) encapsulated into liposomes loaded within the embryo and a near-by array of micro-magnets. In addition, we find that mesoderm constriction and folding leads to the stretching of posterior endoderm cells, which in turn activates the apical accumulation of Myo-II required for its invagination, in a Rho-dependent mechanotransductive process.

Embryonic development has long been considered to be the product of an inductive cascade of developmental gene expression. Our results reveal the additional existence of a rapid self-inducing mechanotransductive cascade of active morphogenetic events, driving the sequence of apical stabilization of Myo-II of the first two invaginations in the morphogenesis of early Drosophila embryos. The cascade consists of the sna-dependent apex pulsations triggering medio-apical stabilization of Myo-II and mesoderm invagination, which in turn mechanically activates the apical stabilization of Myo-II that triggers posterior endoderm invagination.

Results
Micro-magnets to mimic sna-dependent mechanical pulsations. During the early stages of Drosophila embryo development, sna-dependent mesoderm cell apex pulsations initiate at the end of mesoderm cell formation. Using the GFP-tagged transmembrane protein Gilgamesh (Spider–GFP), that labels cell plasma membranes, and Myo-II, which labels the light chain of Myo-II (Myo-II–mCherry), we monitored the change in the surface area of the apex of mesoderm cells and the accumulation of apical Myo-II, respectively. The pulsating phase associated with pulsatile unstable submembranar accumulation of Myo-II, initiated 2 min before the end of cellularization, which was defined as time 0 for all experiments. This phase was observed to last 3 min (Fig. 1b,c). We subsequently observed the initiation of the stable constricting phase, associated with medio-apical stabilization of Myo-II, which drives mesoderm invagination (Fig. 1b,c). Fourier transform analysis of the pulsating phase showed a spectral signature characterized by a mean amplitude of relative surface area of 22 ± 2% with a mean period of 126 ± 22 s (Figs 1d and 2g). In sna RNA interference (RNAi) injected Spider–GFP Sqh–mCherry embryos, both the pulsating and stably constricting phases are absent, no characteristic period emerges from the spectral signature and there was no apical accumulation of Myo-II (Figs 1e–g and 2g). The pulsatile phase was consistently suppressed after injection of the Y27632 inhibitor of Rok that is required for Myo-II activity (Supplementary Fig. 1a).

To test for the triggering of medio-apical stabilization of Myo-II by the mechanical strain generated on mesoderm cells during
snai-dependent pulsations, we developed an innovative approach involving the use of internalized MNPs encapsulated into liposomes and an external array of soft micro-magnets. The soft micro-magnets are magnetized by an external macroscopic magnetic field produced by an electromagnet operated in pulsed mode. When magnetized, the soft micro-magnets produce spatially varying high magnetic field gradients, which serve to apply pulsed forces on superparamagnetic nanoparticles internalized in mesoderm cells (Fig. 1h,i; Supplementary Figs 1b and 2a). By magnetizing the micro-magnets in a pulsed manner, we effectively apply a pulsed force. The resulting dilating-constricting mechanical forces produced within the cells neighbouring the magnets were used for in vivo quantitative mimicking of apex pulsations.

Specifically, the design of the micro-magnets was guided by the fact that in the wild type (WT), both pulsating cells and first neighbouring the magnets were used for dilating-constricting mechanical forces produced within the cells.
constricting cells are found along the ventral midline of the mesoderm. A linear array of iron micro-magnets, approximately cubic in shape and of edge length 10 μm, separated by a distance of 20 μm, was embedded just below the surface of a transparent polydimethylsiloxane (PDMS) membrane of total thickness close to 50 μm (Fig. 1i; Supplementary Fig. 2b; for fabrication details see the Methods section). Liposomes containing superparamagnetic particles, hereafter referred to as ultra-magnetic liposomes (UMLs, see the Methods section), were injected into the ventral cells of embryos (Fig. 1h and Supplementary Fig. 1b). Once injected embryos were positioned ventrally at a distance of 5 ± 2 μm from the micro-magnets (Fig. 1i), a macroscopic homogenous external magnetic field of 180 mT, produced by an electromagnet operated in pulsed ‘60 s on–60 s off’ mode, was applied in the plane of the micro-magnet array, perpendicular to the linear axis of the array (Supplementary Fig. 2b). The superparamagnetic particles are subjected to forces induced by the total magnetic field gradient $\nabla (B_{\text{tot}})$ on the
length scale of the micro-magnet size, that is, of the order of tens of microns (Fig. 1i; Supplementary Fig. 2a–d). These forces are proportional to and in the direction of the gradient of the magnetic field amplitude (that is, of the strongest magnetic field $B_{\text{sat}}$ amplitude: $F = m_{\text{sat}} \nabla (B_{\text{sat}})$)35, if magnetic moments $m_{\text{sat}}$ are saturated by the magnetic field, which, as a good approximation, is the case everywhere at $z = 5 \pm 2 \mu m$ from the surface of the magnets (Fig. 1i; see Methods).

The external magnetic field produced by the electromagnet $B_{\text{ext}}$ serves to magnetize the soft iron micro-magnets, thus generating a magnetic dipole hereafter referred to as the induced field $B_{\text{ind}}$ around each micro-magnet, which leads to the local magnetic field gradients (Fig. 1i; Supplementary Fig. 2c). The existence of this induced local magnetic field gradient was experimentally demonstrated with a scanning hall probe microscope (SHPM), the mean field profile being in quantitative agreement with simulations (Fig. 2a, up; Supplementary Fig. 2d; Methods).

The induced magnetic field $B_{\text{ind}}$ above and below the micro-magnets opposes to the external magnetic field $B_{\text{ext}}$ and leads to a reduced total magnetic field amplitude of $B_{\text{tot}} = B_{\text{ext}} + B_{\text{ind}}$ compared with the amplitude $B_{\text{ext}}$ (Fig. 1i; Supplementary Fig. 2c). Therefore, the presence of the pillar generates local gradients of the total magnetic field amplitude $B_{\text{tot}}$ with smaller values directly above and below the pillar compared with above and below the edges of the pillar (Fig. 1i; Supplementary Fig. 2c, down). The magnetic forces applied to superparamagnetic UML-loaded cells being proportional to, and aligned along the magnetic field gradient, this generates magnetic forces applied on the cells along $y$, oriented from above the pillar to away from the pillar (Fig. 1i). The existence of such repulsive forces is predicted by simulations (Fig. 2a, down; Supplementary Fig. 2e,f). To validate our force simulations, and to demonstrate the existence of repulsive forces in certain regions, we carried out force measurements on a test object (1.5 μm sphere of polystyrene containing superparamagnetic nanoparticles) scanned above the array of micro-magnets, in the presence of an in-plane field (see Methods). Repulsive forces were indeed measured, and good agreement was achieved with simulations (Supplementary Fig. 2g). If the UMLs were free to move, they would follow the force lines and thus move towards the micro-magnets (Figs 1i and 2a; Supplementary Fig. 2f). UMLs cannot aggregate or move inside a cell, since they are 200 nm in diameter, and thus trapped within the smaller mesh of the cytoskeleton which generally is of the order of 70 nm (see Methods)36. The magnetic force acting on the cytoskeleton produced by the UML–micro-magnet interactions is thus transmitted by the cytoskeleton in which UMLs are trapped, to the overall cell. Such forces dominate on the most apical part of the cell (see Methods), and thus dilute UML-loaded cells apex below and in front of the pillars, while they compress those below and adjacent to the pillars along the dorsoventral $y$ axis, forces being negligible along the antero-posterior $x$-axis (Figs 1i and 2a and Supplementary Fig. 2f). The application of the external magnetic field in pulsed mode, with a period comparable to the $sna$-dependent period of apex pulsations (Fig. 1b–g), thus results in soft constriction and dilation pulsatile forces on mesoderm tissue cells in the vicinity of the UML-loaded cells (Figs 1i and 2a). The apical constriction force applied to one UML in one cell was estimated to be of the order of 2 pN, leading to a force applied to one cell apex of $0.5 \pm 0.2$ nN (Fig. 2a; Supplementary Fig. 2d–g with Methods).

To test our ability to quantitatively mimic $snail$-dependent pulsations in magnetically loaded embryos using a soft micro-magnet array coupled with a pulsating magnetic field, we first studied the ventral midline of $sna$ RNAi $twi$ RNAi-injected Spider–GFP embryos loaded with UML. These lack both active $sna$-dependent pulsations31, and the reactive Fog mecha-motion sensitive pathway required for apical stabilization of Myo-II, downstream of $twi$38. In response to the pulsed localized field gradient forces, we observed the generation of individual cell apex size pulsations in the vicinity of the micro-magnets (Fig. 2b; Supplementary Movies 1–3). These pulsations quantitatively mimicked and rescued the $sna$-dependent pulsations (Fig. 2b–d), with a spectral signature showing a characteristic mean period of $138 \pm 28$ s and a mean amplitude of $17 \pm 4$% (Fig. 2e–g). This signature is comparable to the values of $126 \pm 22$ s and $22 \pm 2$%, which are characteristic of $sna$-dependent pulsations (Figs 1d and 2g).

These observations demonstrate the quantitative mimicking of individual cell $snail$-dependent mechanical pulsations within the multicellular mesoderm of the developing embryo, in a fully biochemically independent process, by the application of pulsating magnetic field gradient forces on UML-injected epithelia.

**Mechanical induction of medio-apical stabilization of Myo-II.** We then applied pulsating localized magnetic field gradient forces to Spider–GFP, $Myo$–$m$Cherry embryos injected with both blue fluorescent liposomes (Supplementary Fig. 3b) and $sna$ RNAi only, starting 2 min before the end of cellularization, so as to mimic $sna$-dependent pulsations in the WT.

Focusing on $Myo$–$m$Cherry, we found a medio-apical accumulation of $Myo$-II, with stable accumulation initiated after a time of 50 s (Fig. 3a,b; Supplementary Fig. 3a–c), rescued at 74% of the WT value at 200 s (Fig. 3c; Supplementary Fig. 4). Such $Myo$-II accumulation does not occur in non-pulsating $sna$ RNAi-injected embryos, nor in embryos injected with UML and $sna$ RNAi but not treated to magnetic stimulation (Figs 1e,f; Supplementary Fig. 5; Methods), nor in mechanically stimulated $sna$ RNAi $twi$ RNAi-injected embryos (Fig. 2b; Supplementary Fig. 6a,b). Interestingly, one pulse was sufficient to trigger the response in $sna$ RNAi-injected embryos (Supplementary Fig. 6c).

Looking at both the apical submembrane and junctional planes, we observed that magnetic stimulation led to the rescue of stable medio-apical localization of $Myo$-II, beginning 50 s after the initiation of stimulation, both in junctions (co-localization with $Gilgamesh$ at cell–cell contact in orange) and medio-apically (in red; Fig. 3e), followed by stable collective constriction (Supplementary Movies 4–6). Both sub-apical junctional and medio-apical submembranar localization of $Myo$-II mechanical rescue was also observed in $halo$ $sna$ $Myo$-II–GFP mutants (Fig. 3f), in which the $halo$ phenotype allowed us to select homozygous mutants of $sna$ a priori during cellularization and before gastrulation (Supplementary Fig. 1c)28.

Focusing on Spider–GFP, we found that magnetic stimulation quantitatively mimicked the first phase of $sna$-dependent pulsations (Fig. 3b). In contrast to $sna$ RNAi $twi$ RNAi-injected embryos, in which stable collective constriction and invagination were not triggered (Fig. 2d; Supplementary Movies 7 and 8; Supplementary Fig. 6a,b), the pulsatile phase of $150 \pm 12$ s, comparable to the $170 \pm 10$ s of the WT, was followed by stable apical constriction of cells, with the constriction rate rescued at nearly 50% of the WT (Fig. 3b,d). Associated medio-apical accumulation of $Myo$-II preceded stable apical constriction of mesoderm cells by $100 \pm 12$ s (Fig. 3b), a value also consistent with the $110 \pm 10$ s separating the two events in the WT (Fig. 1c)28.

Thus, the magnetically induced quantitative mimicking of the $sna$-dependent pulsations that rescue medio-apical accumulation of $Myo$-II in $sna$ embryos, also rescued stable medio-apical constriction and mesoderm invagination from both the $sna$ RNAi and $halo$ $sna$, $sna$-defective phenotype (Fig. 4a; Supplementary Movies 4–6 and 9–11; Supplementary Figs 3a–c,e, 7).
We then checked, in time lapse in vivo, that a simple mechanical indent of 20 cells by a micromanipulated needle rescued both the apical accumulation of Myo-II and mesoderm cell constriction in snail mutants, initiating at the location of the indent, then propagating all along the mesoderm at a velocity of 0.29 μm s⁻¹ (Fig. 4b; Supplementary Movie 12). This was confirmed in halo snail mutants (Supplementary Fig. 3d), with a lower efficiency than in the snail indent and in the halo snail magnetically rescued embryos (Fig. 4b Fig. 3f), presumably reflecting the cumulative perturbing effect of the halo genetic mutation and mechanical indent, two non-physiological parameters. Combined with the 150 s of the morphologic contractile response to magnetic deformation (Fig. 3b), one predicts a range of passive deformation by a minimum of 20 constriciting cells initially deformed by the indent able to activate constriction of 0.29 × 150 = 45.5 μm. This represents 8.7 neighbouring cells of 5 μm size. In a linear approximation, this predicts that three cells but not two would be able to deform enough the neighbouring cells to trigger the stable constriction of the nearest-neighbour cells.

Such rapid autonomous propagation of active constrictions confirms that apical constrictions are mechanically activated in response to physiologically constricting neighbouring domains, rescued from the snail mutant phenotype.

These results show that the mechanical cues generated by cell pulsations allowed by snail expression are sufficient to trigger the coordinated apical accumulation of Myo-II leading to coordinated apical constriction and to mesoderm invagination.

**Medio-apical Rok/Myo-II Fog-dependent mechanical stabilisation.** The rescue of apical accumulation of Myo-II leading to constriction and mesoderm invagination by an indent of snail mutants with a needle was suggested to be induced by the mechanical blocking of Fog endocytosis due to membrane strain-induced effects, thereby leading to the over-activation of the downstream of the Fog signalling pathway. Indeed, rescue was found to depend on the expression of the Fog protein18. Consistent with the fact that Fog is expressed downstream of Twist37, we did not rescue apical accumulation of Myo-II, nor collective

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**Figure 3 | Rescue of medio-apical Myo-II in snail-defective embryos by mechanical pulses.** (a) Dynamical imaging of apex size (Spider, Gilgamesh in green) and accumulation of Myo-II (1 μm below, merged in red) in snail RNAi-stimulated embryos in passively deformed and stimulatively sno RNAi-mimicking snail-dependent pulsations in mesoderm cells. Representative of n = 3 of the N = 5 embryos (P = 0.015, exact Fisher test). (b) Quantification on one representative cell of a. (c) Quantification of the apical accumulation of Myo-II at 200 s compared with 0 s in the WT, snail RNAi and snail RNAi magnetically stimulated. N = 6 embryos for each condition (P < 10⁻³ for both snail RNAi versus WT and sno RNAi magnetically stimulated). (d) Constriction rate in WT (N = 6), versus snail RNAi-injected (N = 6, P = 0.004), snail RNAi twi RNAi-injected (N = 6, P = 0.02), twi RNAi (P = 0.002) and sno RNAi magnetically stimulated (N = 4, P = 0.03) embryos. Scale bars, 5 μm. P values are calculated from the Mann-Whitney test, and error bars are s.e. (e) Spider (Gilgamesh, in green) and Myo-II (in red) cortical submembrane and junctional apical locations in the mesoderm cells of WT, snail RNAi and sno RNAi magnetically stimulated embryos, at t = 100s. (f) Myo-II (in green) cortical submembrane and junctional apical locations in the mesoderm cells of WT, halo snail and halo snail magnetically stimulated embryos, at t = 100s. Representative of n = 8, n = 8 (P < 10⁻³, exact Fisher test) and n = 6 embryos (P = 7 × 10⁻³) respectively, of the N = 8 embryos observed in each case.
Constriction nor mesoderm invagination in twiRNAi snaRNAi Myo-GFP embryos submitted to micro-magnet induced mechanical stimulation mimicking sna-dependent pulsations (Fig. 2b,d). Only a partial apical constriction deformation of several mesoderm cells was observed due to local pulsatile magnetic stimulations, thus mesoderm invagination did not occur (Supplementary Fig. 6a,b). We then submitted sna-defective embryos injected with Fog RNAi, to the magnetic stimulation that generated sna-dependent mimetic apex pulsations in sna RNAi-injected Myo-GFP UML-loaded embryos. These embryos showed no detectable rescue of apical accumulation of Myo-II (increase of apical Myo-II to 1.04 ± 0.14 compared with before stimulation), nor mesoderm cell constriction nor mesoderm invagination observed in both the WT and the stimulated sna RNAi-injected embryos (Fig. 5a–c; Supplementary Movie 13). Similar results were consistently observed with Rok, for which medio-apical relocalization is required for Myo-II medio-apical relocalization (Supplementary Fig. 8a–c; Supplementary Movies 14–17).

These results show that the Fog protein, which is downstream of Twist, is required in the sna-dependent mechanically induced medio-apical polarization of Rok that leads to mechanically triggered medio-apical accumulation of Myo-II, apical constriction and mesoderm invagination.

**Mesoderm apex constriction triggers endoderm apical Myo-II.**

Apical accumulation of Myo-II leading to epithelial invagination and folding was first identified in the posterior endoderm formation of gastrulating *Drosophila* embryos. Strikingly, while Sna and Fog were found to be required for apical accumulation of Myo-II and proper invagination of the mesoderm, Fog, but not Sna, is expressed in the posterior endoderm. On the other hand, stable apical accumulation of Myo-II in the posterior endoderm initiates at 223 ± 126 s, after contractile stabilization in the mesoderm at 99 ± 48 s (Fig. 6a,e). The accumulation is subsequently reinforced during mesoderm folding until mid-invagination, before the initiation of the germ-band extension morphogenetic movement of convergent extension (Fig. 6a,e). This correlates with the initiation of the posterior endoderm cell stretching observed in response to constriction and folding of ventral mesodermal cells. We thus tested the hypothesis that the mechanical strain dilations developed by mesoderm cell constriction and folding on the posterior endoderm epithelium activates the apical stabilization of Myo-II, in a Fog-dependent process.

We first characterized apical accumulation of Myo-II in posterior endoderm cells in non-invaginating halo sna mutants. In contrast to the WT, we found no ventral stretching of posterior endodermal cells as observed directly from the position of posterior pole germ cells, and confirmed in the kymographs of posterior endodermal cells (Fig. 6b,d). In contrast to the WT, no apical accumulation of Myo-II in endoderm cells was observed up to t = 204 s (Fig. 6e), that is, during the equivalent time of the WT first stochastic and second collective constriction phases. We
rescued the mechanical effect of stretching of posterior endoderm cells by mesoderm invagination in the halo sna mutant, by attracting UML-loaded posterior mesoderm cells with an array of hard (that is, permanently magnetized) micro-magnets of size 20 μm deposited onto a patterned Si substrate, attached to a micro-manipulator, 6 min before the end of cellularization (Fig. 6c,d). We observed a rescue of the apical accumulation of Myo-II in both the ventral and dorsal endoderm beginning at $t = 12$ s, namely, 4 min after the initiation of stretching, and this was re-enforced between $t = 204$ s until $t = 348$ s (Fig. 6e). The same dynamics of rescue were observed when the stretching was generated from the dorsal side of the embryo, that is, independently of any rescue of mesoderm invagination by magnetic manipulation in the mesoderm, thereby excluding any mechanical-independent developmental rate effect associated to mesoderm formation per se (Supplementary Fig. 9). Given the nearly 3 min delay between the initiation of the magnetic force in the mesoderm and the observation of deformation in the posterior endoderm (Fig. 6d), the timescale of Myo-II apical stabilization rescue is nearly 1 min. This is consistent with the observation of a response initiating nearly 1 min after the first appearance of apical myosin in the mesoderm of WT (Fig. 6a,e), as well as with the 1 min timescale of the Myo-II response to magnetically induced and the < 2 min timescale indent induced, mechanical strains observed in the mesoderm (Figs 3b and 4b, right, respectively). Such rescue was impaired in embryos previously injected with Fog RNAi (Fig. 6e).

Germ-band extension (GBE), which compresses posterior pole cells 8 min after they have been stretched by mesoderm invagination (Fig. 6a), could subsequently also participate in the posterior endoderm apical accumulation of Myo-II and invagination in a sna-independent process, after stretching due to mesoderm invagination. We thus tested the role of the sna-dependent mesoderm invagination induced mechanical activation of apical relocation of Myo-II, in endoderm invagination under bcd, nos, tor germ-band defective conditions$^{13}$. In sna and bcd, nos, tor RNAi-injected embryos, lacking both mesoderm invagination and germ-band extension, no invagination was observed in the posterior endoderm (Fig. 6f). The magnetically induced mimicking of posterior endoderm cell stretching by sna-dependent mesoderm invagination rescued shallow cup formation, that is, the initiation of endoderm invagination$^{41}$, in these embryos (Fig. 6f). Consistently, bcd, nos, tor RNAi-injected embryos showed shallow cup formation initiation, associated with sna-dependent mesoderm invagination (Fig. 6f; Supplementary Movies 18 and 19; Supplementary Fig. 10e). Note that here the initiation of endoderm invagination is less pronounced than in the WT, possibly due to the lack of an additional germ-band extension buckling driving force in the bcd, nos, tor RNAi-injected embryos. Similarly, bcd, nos, tsl germ-band defective mutant embryos are also known to initiate posterior endoderm invagination through a shallow cup formation$^{44}$, as observed here with posterior folding and the cup shaped yolk–epithelium interface, respectively—two red arrows in Fig. 6f. Finally, sna-defective embryos showed apical accumulation of Myo-II and the initiation of invagination associated to bcd, nos, tor-dependent germ-band extension at later stages after 600 s (Fig. 7).

These experiments demonstrate the Fog-dependent mechanical induction of posterior endoderm apical stabilization of Myo-II and invagination initiation triggered by the mechanical strains developed by mesoderm constriction and invagination, which is subsequently reinforced by morphogenetic movement of GBE.

**Discussion**

The search for the biochemical principles regulating epithelial cell folding leading to tissue invagination is a long-standing challenge in developmental biology$^{30}$. Specific interest has been focused on the endoderm and mesoderm primary invaginations of gastrulation, which correspond to the very first morphogenetic movements of embryogenesis in many distinct species$^{45}$. One of the best genetically characterized morphogenetic...
movements is mesoderm invagination of early Drosophila embryos, which is regulated by the interplay between the two pathways established by the expression of the first zygotically expressed ventral genes snail and twist. However, while the expression of snail is known to be required for the twist-dependent stable medio-apical accumulation of Rok and Myo-II, mesoderm cell constrictions and invagination, the underlying mechanism of the interplay between the two pathways remains poorly understood.26,31,38,46,47. Because snail expression leads to cell apex size pulsations through a process that remains to be precisely characterized, thereby introducing mechanical strains in mesoderm cells, it was experimentally (by using a simple but non-physiological indent mechanical perturbation) and theoretically (by using simulations) suggested that snail could activate the Rho pathway leading to apical accumulation of Myo-II, through a downstream of twist, Fog-dependent mechanotransductive process.18,48. This hypothesis was in line with the finding of a twist-dependent increase of probability of having stable apex constriction in the cells neighbouring already constricting cells.49.

The results presented above demonstrate that quantitative mimicking of snail-dependent pulsations in the absence of snail expression, leads to (i) the rescue of medio-apical stabilization of Rok, (ii) apical and medio-apical stabilization of Myo-II, (iii) coordinated stable apical constriction and mesoderm invagination. All these processes, absent in snail mutants, are rescued on the timescale of a few minutes. In addition, we show that the rescue of these WT phenotypes from the snail phenotypes is Fog-dependent. This demonstrates the existence of a Fog-specific
mechanotransductive activation of medio-apical polarization of Rok, leading to apical stabilization of Myo-II by the mechanical strains resulting from sna-dependent pulsations in the mesodermal cells of Drosophila embryos, on the timescale of a few minutes.

Interestingly, the fact that the rate of constriction rescue is nearly two times lower in the absence of sna than the constriction rate of the WT, and has a lower intensity of Myo-II accumulation at the timescale of invagination initiation (>300 s), possibly reflects the lack of one of the receptors of Fog (that is, Mist) in sna mutants (Fig. 3b,c). Nevertheless, these less intense responses to mechanically induced pulsations led to mesoderm invagination rescue in sna mutants, with similar dynamics to that of the WT, showing that sna-dependent pulsations are sufficient to trigger the medio-apical stabilization of Myo-II required for mesoderm invagination.

Invagination specifically requires the coordinated constriction of apexes. Coordination of individual cell behaviour can be generated in response to a common external biochemical morphogenetic field, such as the expression of the Fog secreted signal protein. However, ectopic expression of Fog in sna mutants, as well as Fog overexpression in the mesoderm of sna-defective embryos, do not lead to apical accumulation of Myo-II, nor apical constriction nor mesoderm invagination. Coordination of individual cell behaviour can be achieved if the behaviour of a given cell activates the same behaviour in its neighbouring cells though cell–cell interactions, thereby leading to correlated collective effects. In the specific case dealt with here, the individual cell behaviour is apical constriction, and is thus mechanical in nature. It was suggested that the mechanical activation of cell constriction by a neighbouring constricting cell can drive collective and coordinated constriction. Here we propose that sna-dependent pulsations represent a phase of unstable autonomous cell apex constrictions generating cell–cell mechanical interactions, which trigger a transition to a phase of coordinated apical constriction. Such a transition emerges from a Fog-dependent mechanotransductive activation of the stable constriction of cells mechanically deformed by their neighbouring pulsating constricting cells.

Invagination of both the mesoderm and posterior endoderm is Fog-dependent and induced by a process of apical accumulation of Myo-II. Strikingly, we show here that the initiation of both apical Myo-II accumulation in the mesoderm and invagination in the posterior endoderm require the expression of sna (Fig. 6e), while the posterior endoderm does not express sna. However, we found that sna is not biochemically required in itself for the medio-apical accumulation of Myo-II leading to mesoderm invagination, rather the mechanical strains developed by sna-induced pulsations are needed. Similarly, we show that the triggering of apical accumulation of Myo-II and posterior endoderm invagination is induced by the mechanical strain dilations developed by the constriction of mesoderm cells, on the timescale of a few minutes. We also establish the conservation of the Fog-dependent mechanotransductive induction of Myo-II apical accumulation by mesoderm invagination, in the posterior endoderm. Thus, the fact that sna is required to mechanically trigger apical accumulation of Myo-II and mesoderm invagination implies that sna expression is also required to initiate posterior endoderm invagination, even though its expression is not required in posterior endoderm cells. As a consequence, mesodermal sna interacts with posterior endodermal Fog, over the distance of five cells separating the two tissues, through mechanical cues, known to propagate rapidly and over long distances across materials. In addition, we find that mechanical induction of Myo-II apical accumulation and invagination in the posterior endoderm is subsequently reinforced by the morphogenetic movement of GBE.

Embryonic development has long been considered as the product of a cascade of developmental gene expression, with characteristic timescales ranging from 10 min to hours for each individual step of the cascade. Here we demonstrate that the Fog-dependent mechanotransductive induction of medio-apical polarization of Myo-II regulates a self-inductive non-transcriptional rapid cascade of active morphogenetic movements at the shorter timescale of minutes. Sna-dependent mesoderm apex pulsations trigger the Fog-dependent medio-apical stabilization of Myo-II that activates the coordinated constriction of apexes generating mesoderm folding and invagination within a few minutes. Such coordinated constriction and mesoderm invagination in turn activates the Fog-dependent apical accumulation of Myo-II initiating the next morphogenetic movement of posterior endoderm invagination, which also occurs within a few minutes (Fig. 8). Snail and Fog expression being the result of the biochemical cascade of differentiation, we thus show the coexistence at different timescales, and interactions between a slower differentiation gene expression biochemical cascade and a rapid self-inductive mechanotransductive cascade. Here we thus reveal that, during and patterned by the expression of developmental differentiation genes, mechano-protein mechanotransductive interplay drives a rapid self-inducing cascade of morphogenetic movements in early Drosophila gastrulation.

**Methods**

**Fly strains.** Oregon R was taken as the WT. Male and female of any age were used. Myo-II-GFP (sqa-GFP) was provided by Roger Karres, Myo-II-mCherry N2/rok-GFP and N2/Gilgamesh–GFP (sqa–mCherry spider–GFP), snail-GFP (small Hill05/CyO, sqh–sqa–GFP) were provided by Adam Martin, sna-GFP (small Hill05/CyO, sqh–sqa–GFP) were produced by crossing holo small Hill05/CyO, sqh–sqa–GFP with small Hill05/CyO, rok-GFP (w; Ubi-Rok-GFP; et w; Ubi-Rok-GFP/TM6) were produced by Yohanns Bollache. Stocks were maintained at room temperature (22°C), and experiments and observations have been done at 24°C.

**RNAi.** Three RNAi sequences a gene were designed and produced by Eurogentec: Sna RNAi 1 (1,481–1,503 S 5′: EACCUAUAAGUUACAGGAT dTdT; AS 3′: ACGCUUUAAUACAGGAT dTdT), Sna RNAi 2 (340–362 S 5′: GGAACCCAGAACGCUUACUACU dTdT; AS 3′: AUAGUGUACUGUAGUACUUACUACU dTdT), Sna RNAi 3 (972–994 S 5′: CCGAGUGUAAUCAGGAGAA dTdT; AS 3′: AUAGUCACGUUUCGGUUCC dTdT), Sna RNAi 4 (1,005–1,027 S 5′: GCACCAGCAACAGAUCUAU dTdT; AS 3′: GCACCAGCAACAGAUCUAU dTdT), Sna RNAi 5 (2,061–2,083 S 5′: GUCACGCUUUACUAUAUA dTdT; AS 3′: UUAUAUAUGGAAAGCGUGAC dTdT).

**Figure 7 | Distribution of Myo-II-GFP in halo sna at germ-band extension initiation.** Representative of the N = 6 WT embryos and the N = 6 halo sna embryos (P = 2 × 10⁻³, exact Fisher test). Halo sna invaginate their posterior endoderm (later stages, not shown), as do sna mutants. Scale bars, 5 μm.
**Figure 8 | Meso-endoderm mechanotransductive cascade of Myo-II apical stabilization.** 

(a) Sna-dependent stochastic fluctuation of mesoderm apex size (b) activates the medio-apical accumulation of Myo-II in a Fog-dependent mechanotransductive process leading to coordinated collective constriction of the cell apex (initiating nearly 130 s after the initiation of pulsations), then mesoderm invagination. (c) In turn, mesoderm constriction generates a strain deformation of posterior endoderm cells (blue array) that activates the apical stabilization of Myo-II required for endoderm invagination nearly 80 s after mesoderm collective constriction initiation, that is subsequently reinforced during mesoderm invagination, in a Fog-dependent mechanotransductive pathway. Adapted from the figures of morphogenetic movements of Sweeton et al., Development, 1991 and Nüslein-Volhard & Wieschaus, Drosophila a practical approach, 1896.

CGGUAACCGACUAAUGU dTdT; AS 5': ACUAAUGUGUUGCUGUACCG dTdT.

Fog RNAi: (98–12,520 S 5': CUGGCGCAGAUGUAGUAU dTdT; AS 5': AUCU AAUACCUGUGGCG dTdT, 7054–7076 S 5': CUGAACACCACAU AAUA, 7077–7100 S 5': UAAUUGUCGGGUGUACCG dTdT, 7101–7134 S 5': CGUGGGAGGGUUGUAUCCAA dTdT; AS 5': UUGAAUAUACAAUCUCAC ACG dTdT).

The sequences were designed to not interfere between themselves, and with myo-II, rhogef2, EBI, ctu, Tafq, trafj, Jupiter, arm, src2A, src42B, shg, khc, khc. Bicoid-RNAi (3,592–3,614 S 5': UUGAAAUCAACUUCAC dTdT).

Bicoid-RNAi was not currently available (in the community) to test for Torso protein.

Twi RNAi induced the specific decrease of 46% of Twist expression as a mean value of the overall stage 5 (with no decrease of Bicoid expression as a control; Supplementary Fig. 10b), as mean values characteristic of the overall stage 5, which together was sufficient to lead to the specific inhibition of germ-band extension at mid-stage 6 (Supplementary Fig. 10c). Note that the Torso antibody was not currently available (in the community) to test for Torso protein.

Thus, RNAi treatments let to a specific decrease of the order of 70% of the signalling protein Fog, of 40% of the targeted proteins Snail and Twist, and of the order of 30% on the Bicoid and Nanos proteins having a combining effect in the production of germ-band extension, during stage 5. Such quantitative gradation is consistent with the inversely related to the developmental time at which the RNAs were produced and supporting the specific inhibition of germ-band extension at mid-stage 6 (Supplementary Fig. 10c).

**Embryo preparation.** Egg collections were performed on agar cups covered with grape-juice plates. Flies were allowed to lay eggs for 2–4 h at 25°C before the plate was removed. During preparation, embryos were never compressed to avoid any process of mechanotransduction before controlled mechanical stimulation. In all experiments, Drosophila embryos were dechorionated with 50% bleach until the ‘antenna’ disappear (nearly 40 s), then washed with water. Embryos were then glued on a coverslip covered with glue (double-sided tape soaked in heptane), and then left to dry. In case of injection, embryos were previously dried 7 min in the presence of salt. A few droplets of Halocarbon oil 27 were placed on the embryo to prevent drying and to allow observation. After injection, an area of 1 cm diameter including the somite was cut with a razor blade, cutted upside down and glued on the piezo-electric system to observe the ventral side in reflection, and to allow controlled relative movements of the embryo with respect to the observing system.
to the magnetic pattern with 100 nm resolution. For indenter experiments, the coverslip with the glued embryo was put on the microscope stage, and the indenting needle was attached to the piezo-electric device.

Ultra-magnetic nanoparticle preparation. The synthesis of MNPs is described first. The aqueous suspension of MNPs was prepared using alkaline co-precipitation of FeCl$_3$ (0.9 mol) and FeCl$_2$ (1.5 mol) salts, according to Massart’s procedure 51. Superparamagnetic maghemite grains (γ-Fe$_2$O$_3$) were obtained by oxidizing 1.3 mol of magnetite with 1.3 mol of iron nitrate (boiling solution). After magnetic decantation, 21 of distilled water and 360 ml of HNO$_3$ 20% were added to the solution and the mixture was stirred for 10 min. Prepared magnetite nanoparticles were washed several times with acetone (3 × 1) and ether (2 × 500 ml) and suspended in water. Size sorting was performed by adding HNO$_3$ (0.45 M) to the suspension followed by magnetic decantation. This operation was repeated with the deposit until suitable particle size was obtained. Sodium citrate (n$_0$/n$_1$ = 0.13) was added to the nanoparticles and the mixture was heated to 80°C for 30 min to promote absorption of citrate anions on their surface. Citrated nanoparticles were precipitated in acetone and suspended in water. The volume fraction and average size of the magnetite grains were determined by fitting the magnetization curve of nanoparticles using the Langevin function. Particles of 9 nm diameter (n.d. = 0.35, volume fraction of nanoparticles in the suspension $\phi = 1.9\%$, specific susceptibility $\gamma/\phi$ of 15.5) were obtained. For UML preparation, the aqueous medium was removed using a Macroset centrifugal device 30 kDa (PALL) and nanoparticles were suspended again in a buffer (0.108 M NaCl, 0.02 M sodium citrate and 0.01 M HEPES, pH 7.4).

The preparation of UMLs is described second. Solutions of 1,2-dipalmitoyl-sn-glycerol-3-phosphocholine (DPPC), 1,2-distearoyl-sn-glycerol-3-phosphocholine (DSPC), 1,2-distearoyl-sn-glycerol-3-phosphoethanolamine-n-[carboxy(polyethylene glycol)]2000 (ammonium salt) (DSPE-PEG2000) and L-sphingosine, sphingomyelin- and cholesterol-containing liposomes were used to produce liposomes. The procedure was performed at constant pressure, with the visualization of UML labelled with fluorescent dye. The magnetization of the Fe micro-magnets was assumed to be given by the demagnetizing susceptibility so that $M = H_{demag} N$, where $N$ is the demagnetizing field factor of the micro-magnet, which is taken to be 1/3. This gives $M = 0.18 \times 3 - 0.54 T$. The applied field used in the calculation was that produced by the bulk magnet used to magnetize the micro-magnets in the SHPM. The total field of the field profiles is comparable, though the peak-to-peak value is higher in the measurement (34 mT) than in the simulation (18 mT; Fig. 2a; Supplementary Fig. 2d). This difference may be attributed to a number of factors, including the non-cubic shape of the real micro-magnets, an error in the estimate of the sample-to-probe distance, and a contribution from the side-wall deposits, not included in the simulations. Nevertheless, the agreement is considered sufficient for support force simulations shown in Fig. 2a; Supplementary Fig. 2e, detailed below.

Quantification of magnetic forces. First, evaluation of forces produced on one UML. The energy of interaction of the UML, with an induced magnetic moment $m_{UML}$ and the total magnetic field, $B_{tot}$ is: $E = - m_{UML} B_{tot}$. (ref. 35). The magnetic force $F = - \nabla E = - V (m_{UML} B_{tot})$. The magnetic moments of the superparamagnetic nanoparticles in the UML are aligned along $B_{tot}$ so: $V = \sum_{i=1}^{n} \frac{m_{i} m_{tot}}{C_{1} B_{tot}}$, where $m_{i}$ and $B_{tot}$ are now scalar. UMLs at $z = \pm 2\mu m$ above the micro-magnet experience $B_{tot} \sim 140 mT$ (the value is reduced compared to the external field value, because of the opposing induced field) and the nanoparticles are almost saturated. The Supplementary Fig. 1 shows the magnetization curve of the UMLs in ref. 34). Thus, $m_{UML}$ can be considered as almost constant, so that the force applied to the UML is proportional to the gradient of the amplitude of the magnetic field, and is given by: $F = m_{UML} \nabla B_{tot}$. Note that because of the $8 \times 10^{-10}$ concentration of UMLs found in cells after injection (see Evaluation of the number of UML inside mesoderm cells paragraph below), the mean distance between UMLs within cells is thus of the order of 10 times larger than their size (200 nm), so that attractive interactions between magnetized UMLs are very small compared with the force of interaction between the UMLs and the micro-magnets, given the $\left(\frac{d\gamma}{d} \right)^2$ decay of the interactions between magnetic moments of magnetic systems of size at a distance $r$ (ref. 35). What is more, UMLs cannot aggregate, since they are trapped within the mesh of the cytoskeleton which is generally of the order of 70 nm (ref. 36). Effectively, no apical concentration of UMLs is specifically observed under the magnets after the application of the magnetic field pulses, in the sns twi-defective mutants in which the observation is made easier by the fact that invagination cannot be activated, demonstrating that the UMLs are trapped in the cytoskeleton (Supplementary Fig. 6b, t = 435 s). In addition, as confirmed by Supplementary Fig. 1b, UMLs are inside cells, and the UMLs that would remain present out of the cell but within the embryo into the yolk will not significantly contribute to the force applied to the cells, because they are positioned at $>20\mu m$ from the yolk. In that respect, it is obvious that the specific range of the applied magnetic force (magnetic forces decay strongly with the distance to the magnet, and practically, only UMLs in the 5μm most apical part of the cell feel significant force, see the first section in the Methods section ‘Quantification of magnetic forces’ and paragraph below associated Supplementary Fig. 2d). In addition, there is no extracellular matrix around the cells into the yolk, and more generally at a stage of development, to trap the UMLs. The magnetic force acting on the cytoskeleton is thus dominated by constricting/dilating forces produced by the UML–micro-magnet interactions, transmitted by the cytoskeleton in which UMLs are trapped, to the overall cell.

Injection of UML and RNAi. An injector (Eppendorf Femtotjet) was fixed on the microscope to inject magnetic particles/or and/or RNAi. A x:y:z piezo-electric system (Princeton Instruments) was mounted on the microscope to control the position of the embryo mounted on a coverslip relative to the injecting needles, the micro-magnets or the indent needle, with 100 nm resolution. For indent experiments, the coordinates of the embryo mounted on a coverslip relative to the injecting needles, the micro-magnet array (Princeton Instruments) was mounted on the microscope to control the position of the embryo mounted on a coverslip relative to the injecting needles, the micro-magnet array (Princeton Instruments). Using a Femtojet Eppendorf injector, approximately 0.05 ml of 0.1 g of UML was injected into embryos along by cell basal domain inside the ventral cells of the embryo from 40 to 20 min before the end of ventral cellularization. Injections were performed at constant pressure, with the visualization of UML labelled with Rhodamine observed by confocal microscopy. Visualization of the Rhodamine labeled stage in the buffer was introduced before sonication at room temperature for 20 min to produce a water-in-oil emulsion. Organic solvent was evaporated with a rotavapor R-210 (Buchi) at 25°C. Liposomes were filtrated through a 450 nm filter until the gel phase disappeared. Liposomes were filtered through a 450 nm filter and purified from nonencapsulated MNPs by magnetic sorting using a bulk magnet (Calamari–Fe – Nd – B 150 × 100 × 20 μm).

Fabrication and simulation of micro-magnets. Deep reactive ion etching was used to produce a linear array of Si pillars of cross section 10 × 10 μm and height 40 μm, separated by a distance of 20 μm, at the surface of a Si substrate. High-rate triode sputtering was then used to deposit a Ta (100 nm)/Fe (10/μm)/Ta (100 nm) tri-layered structure based on the Ta being a capping layer (the Ta layer being thicker to protect the Fe layer from reaction with the substrate and from oxidation, respectively). The thickness of Fe deposited on the sidewalls of the pillars is limited due to the directional nature of sputtering and to shadowing effects from neighbouring pillars. A 50 μm-thick layer of PDMS was spin-coated onto the structure and, after relictuation, the PDMS layer was pulled off the substrate in such a way that the embedded magnet-topped Si pillars break away from the substrate. The resultant structure consists of a PDMS membrane containing a linear array of soft micro-magnets of approximate cubic shape, just below the membrane surface. More details can be found in ref. 62.

A SHPM was used to measure the stray field pattern produced by the micro-magnets 64. The micro-magnet array was positioned on top of a bulk NdFeB magnet, oriented in such a way to produce a field that is comparable to that applied by the electromagnet during bio-experiments (in-plane, oriented perpendicular to the axis of the array of micro-magnets). The Hall probe was mounted at an angle of ~23° with respect to the y axis, for experimental reasons, so that the measured signal had contributions from both the $z$ and $y$ field components ($B_x$ cos23° + $B_y$ sin 23°). The scan was made at a distance of roughly 10 μm above the micro-magnet array. Note that the field component is determined by the orientation of the probe and that it is not possible to measure the component of the field pattern above the micro-magnets using our SHPM.

Nevertheless, the good agreement observed between measured and simulated patterns validates our simulations of all field components. Simulation of the stray magnetic field pattern produced by the micro-magnet array was performed. To estimate the force applied by the micro-magnet array on an UML, we needed to know the magnetic field produced by it. Then, using a SHPM, at a scan height of 10 ± 5 μm, as shown and described above (Fig. 2a). We then simulated the field profile in an analytical approach using Matlab. We considered a linear array of seven cubes of Fe, of edge length 10 μm, separated by 20 μm, magnetized by an in-plane field, applied perpendicular to the axis of the UML. The magnetic moment of the Fe micro-magnets was assumed to be given by the demagnetizing susceptibility so that $M = H_{demag} N$, where $N$ is the demagnetizing field factor of the micro-magnet, which is taken to be 1/3. This gives $M = 0.18 \times 3 - 0.54 T$. The applied field used in the calculation was that produced by the bulk magnet used to magnetize the micro-magnets in the SHPM. The total field of the field profiles is comparable, though the peak-to-peak value is higher in the measurement (34 mT) than in the simulation (18 mT; Fig. 2a; Supplementary Fig. 2d). This difference may be attributed to a number of factors, including the non-cubic shape of the real micro-magnets, an error in the estimate of the sample-to-probe distance, and a contribution from the side-wall deposits, not included in the simulations. Nevertheless, the agreement is considered sufficient for support force simulations shown in Fig. 2a; Supplementary Fig. 2e, detailed below.
To evaluate the force applied on one UML, let us first consider the saturation values of magnetization and moment. The saturation magnetization of the constituent phase of the MNPs, maghemite, is \( M = 3 \times 10^3 \text{ Am}^2 \text{m}^{-1} \). Thus, the saturated moment of a 9 nm magnet is \( M_V = 1.15 \times 10^{-19} \text{ Am}^2 \). Based on magnetic measurements reported in the paper by Bealle et al., the UMLs, of average diameter 200 nm, are filled roughly with roughly 20 vol% of 9 nm diameter MNPs. The saturation magnetization of a given UML is thus 76 mT, and its saturated moment 2.5 \( \times \) \( 10^{-16} \) Am². Calculations show that the magnetic field strength at 5 mm above the micro-magnet array is \( 0.14 \text{ T} \). From the \( M(H) \) curves of UMLs, measurements reported in Supplementary Materials section of ref. 34, \( 0.14 \text{ T} \) gives 90% of saturation. From above the micro-magnet array is line with compression forces of 0.6 nN per cell (60 nN involved in the compression of the UMLs and the MNPs are 80% saturated. Thus, the magnetic moment of a superparamagnetic particle is given by the Langevin function:

\[
m = m_0 \mu (x) = m_0 \mu (x) \frac{H}{H_{sat}}
\]

Using \( m_0 = 1.15 \times 10^{-19} \text{ Am}^2 \), \( \mu (x) = 0.18 \text{ T} \) and \( T = 300 \text{ K} \), \( x = 5 \) so \( m_0 = 80 \% m_\text{sat} \). Thus, we consider that the UMLs and the MNPs are 80% saturated. \( m_\text{sat} = 2 \times 10^{-15} \text{ Am}^2 \) is used in the force calculations. The force on a UML was calculated as a function of its position with respect to the micro-magnets using analytical expressions \( F = m_\text{UML} \nabla (B_m) \) (Matlab; Fig. 2a).

Second, measurement and simulation of force on a test object are described. To validate our force simulations, and to demonstrate the existence of repulsive forces in cells, we carried out force measurements on a test sample scanned above the array of micro-magnets. The test object was a 1.5 mm sphere of polystyrene containing 30 volume % superparamagnetic nanoparticles, supplied by microParticles GmbH. The sphere was stuck to the apex of a commercial atomic force microscope cantilever using a scanning electron microscopy (SEM) glue. The membrane containing the micro-magnet array was placed on the bulk magnet used for SPMH, so as to apply an in-plane field perpendicular to the axis of the linear array of magnets (that is, along y). The superparamagnetic microsphere was then scanned at different heights above the micro-magnet array, and the deflection of the probe along the z axis was measured as a function of position \( z \). The deflection \( x \) (note that the probe was not rotated in resonance mode).

The level of deflection is directly proportional to the z component of the force exerted by the micro-magnets on the microsphere. As can be seen in Supplementary Fig. 2g, the deflection force changes from attractive to repulsive, depending on the position of the sphere. The sphere is repelled when it is above the micro-magnets but attracted when displaced along the y axis, in the direction of the applied field. The force on the test object was calculated using analytical expressions (Matlab). Relatively good agreement is found between the measured and calculated forces. Good agreement was also found between the measured and calculated z derivative of the peak-to-peak force, as a function of \( z \) (Supplementary Fig. 2g). Though the approach used only allows us to measure the z component of the force, not the y component, it nevertheless validates our approach to simulations and demonstrates that both attractive and repulsive forces can be experienced in the geometry used in our bio-experiments, that is, with an in-plane field serving to magnetize the micro-magnets and the UML within the cell.

To quantify the number of UMLs, we measured the intensity of UML fluorescence per cell, subtracted from the background intensity away from the embryo. This led to local patterns of maximum intensity. Taking the smallest intensity maximum as a single UML signal, we evaluated 18 \( \pm \) 7 UML per cell in an optical section 0.5 mm-thick two-dimensional disc (measured in \( N = 6 \) embryos, 494 cells, error bars being s.d., namely, 540 \( \pm \) 210 UMLs in a 15 mm-thick cell of 5 mm apical diameter, or an UML concentration of \( 8 \times 4 \times 10^{-3} \). Third, evaluation of forces produced on one cell is described. The force applied apically to UMLs in one cell to compress its neighbouring cell apex by 20% was thus estimated. The order of magnitude is that the individual UML located in the 5 mm most apical part of the cell are submitted to a mean constraining force of \( 2 \times 10^{-12} \text{ N} \) (Fig. 2a). This order of magnitude is in line with compression forces of 0.6 nN per cell (60 nN involved in the compression of 200 cells). Note that the volume of the 3 mm diameter nucleus approximated as a sphere, that apically excludes UMLs, was evaluated to be \( 9 \mu m^3 \), which remains small compared with the volume of the apical cylinder in which the nucleus is located (\( 5 \mu m \); length: \( 5 \mu m \) so volume = 98 \( \mu m^3 \)), and was thus not taken into consideration in the evaluation.

Conventionally, applied or pulsed homogeneous magnetic field either in the absence of the micro-magnet array (Supplementary Fig. 5) or without injecting UMLs resulted in neither pulsations nor medio-apical accumulation of Myo-II nor apical constriction. This is in agreement with the fact that the UML only feel a field gradient force, and that the micro-magnets are required to locally concentrate the magnetic field so as to produce localized field gradients. Magnetic field gradients also \( a \ priori \) apply a force normal to the tissue, which could be some cases be detected as leading to a maximum of 0.5 mm deformation. However, the mechanical induction of the medio-apical accumulation of Myo-II by normal deformation indentation requires a much higher perpendicular deformation of the tissue, of a minimum of 5 mm (ref. 18), ruling out any effect of the normal forces developed by the micro-magnets in the activation of the medio-apical accumulation of Myo-II.

**Imaging.** A spinning disc (Yokogawa) microscope (IX70 Olympus) coupled to a CoolSnap HQ² camera was used in all experiments using the 405, 491 and 561 nm length wavelengths, as well as transmission. All settings, including the piezo-electric control of z position for stack production (Princeton Instrument), were controlled by the Metamorph software, all assembled by Roper. Images have been taken with a 40 times water objective (Olympus LUMPLN40XW) and a 60 times oil objective (Olympus PLAPL60XO1.42). A 20 times air objective (Olympus UPLSAP020XO.85) was also used for the injections and for some other observations.

Now that the 50 mm-thick PDMS membrane containing the micro-magnets introduces a new optical index medium in between the glass coverslip and the halocarbon oil in which the embryos are immerged, and that defraction is introduced by the presence of the 10 mm iron cubes every 20 mm. This introduces noise in the images that however does not prevent cell resolution analysis. Note also that the embryos are not flattened against the coverslip or PDMS to avoid any mechanically induced interference due to the squeezing. Thus invaginations generated large movements that frequently obliged the operator to re-centre the image to not lose the cell apices. This prevented also the use of \( \times 60 \) magnification during the constriction phase, to not lose the imaging field and mean focalization plan. The use of \( \times 40 \) introduced pixelization in some cases.

**Image analysis.** Apical cell shapes were imaged at junctions, 3 mm in z (in depth) away from the external surface of the test embryos. Individual apical cell areas were measured using FIJI and Free software. Time lapse was realized in vivo every 3 s using spinning disc imaging of the apical domain of the cells of the ventral side of Drosophila embryos labelled with spider–GFP and sph–Cherry. Individual cell tracking were realized afterwards, in silico. Individual cells were manually tracked. The observation started 120 s before the end of ventral cellularization, and ended 300 s after, at the beginning of germ-band extension. Matlab software was used to execute Fourier transform on the dynamics of individual apex cell shape changes. No filtration was applied to generate the Fourier transform. The drift (mean slope of the basal corresponding to stable constriction that adds to pulsations) was removed, to analyse the signal that is associated to pulses only. Mean values and graphs were realized with Kaleidagraph software.

Time-lapse resolution is of the order of 1 min for whole mesoderm imaging. Fluorescent intensity was quantified by subtracting the background intensity (away from the embryo) from the signal measured in cells. Correlation functions between distant pixels on PDMS to avoid any mechanically induced interference due to the squeezing. Thus invaginations generated large movements that frequently obliged the operator to re-centre the image to not lose the cell apices. This prevented also the use of \( \times 60 \) magnification during the constriction phase, to not lose the imaging field and mean focalization plan. The use of \( \times 40 \) introduced pixelization in some cases.

**Drosophila mechanical indents.** A micromanipulated 50 mm diameter needle was approached close to the mesoderm of stage 5. In the specific case of \( sna/Cyogm\) embryos. Embryos having not initiated apical accumulation of Myo-II 2 min after the end of cellular cellularization were gene typed as \( sna \) homorygous. Embryos were systematically indented by 3–7 mm every 4 min at the most posterior 1/3 position of the mesoderm, 2-3 min after the end of cellularisation. After stimulation, the embryo was turned 90° for observation of the ventral side.

**Statistics.** All statistics based on quantitative measurements were evaluated with the non-parametric two-sided Mann–Whitney test, with s.d. error bars being comparable within each group of data, and the variance between groups of data being similar. Error bars are s.d., except in Fig. 3f in which they are s.e. Most of the experiments were reproduced at least six times in the laboratory, which is standard in single embryo in vivo experiments. \( n \) is the number of measurements per experiment. All statistics based on qualitative measurements (on/off results, that is, the presence of an invagination or not) were evaluated with the exact Fisher test.

**Data availability.** The authors declare that all data supporting the findings of this study are available within the article and its Supplementary Information files or from the corresponding author upon reasonable request.

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**Author contributions.** Paul Baggethun.

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**Acknowledgements**

This research was funded by the ANR (grant # 11 BSV5014-01), the AVIESAN Cell Biology, Development and Evolution (CNRS, INSERM, INSTITUT CURIE co-funding of the dedicated Magneto-Spinning disk) and the Labex CeTisPhyBio (grant # 11-LBX-0038), and the Labex LANEF (ANR-10-LABX-51-01). We are very grateful to the following members of Institut Néel: Dominique Givord for his input and guidance throughout the work, Klaus Hasselbach for guidance with magnetic field measurements, Jean-François Motte for preparing the probes for force measurements and Thibaut Devillers for assistance in micro-magnet fabrication, Vasilis Apostolakeas for assistance in illustrations and to Elena Fernandez-Sanchez from Institut Curie for her contribution to western blots analysis. The micro-patternning of Si substrates for micro-magnet preparation was performed in the PTA-Grenoble cleanroom facility.

**Author contributions**

D.M. realized all experiments and B.D. initiated experiments in the posterior endoderm. J.C.-R. realized the westerns. D.L.R., F.D.B. and N.M.D. designed the magnetic tools and D.L.R. fabricated them. G.S. and S.L.D. carried out magnetic field and force measurements, respectively. L.R. carried out field and force simulations. C.M. and A.M. designed and realized the UML. E.F. designed and coordinated the overall work.

**Additional information**

**Supplementary Information** accompanies this paper at http://www.nature.com/naturecommunications

**Competition financial interests:** The authors declare no competing financial interests.

**How to cite this article:** Mitrossilis, D. et al. Mechanotransductive cascade of Myo-II-dependent mesoderm and endoderm invaginations in embryo gastrulation. *Nat. Commun.* **8**, 13883 doi: 10.1038/ncomms13883 (2017).

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