Article

On Implicit Time–Fractal–Fractional Differential Equation

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Abstract: An implicit time–fractal–fractional differential equation involving the Atangana’s fractal–fractional derivative in the sense of Caputo with the Mittag–Leffler law type kernel is studied. Using the Banach fixed point theorem, the well-posedness of the solution is proved. We show that the solution exhibits an exponential growth bound, and, consequently, the long-time (asymptotic) property of the solution. We also give examples to illustrate our problem.

Keywords: well-posedness; exponential growth bound; fractal–fractional operators; Mittag–Leffler type kernel

MSC: 26A33; 28A80; 34A08; 34A09

1. Introduction

The fractal–fractional differential equation is a link between the fractal and fractional differential equations. Fractal and fractional differential equations are known for modeling complex physical processes and phenomena, particularly irregular systems with memory. Although fractional equations are renowned for representing systems with long-term memory and long-range interactivity, fractal calculus, conversely, is immensely effective in working with occurrence in stratified or porous media. That is, fractal–fractional differential operator models physical phenomena and real-world activities that exhibit or display fractional behaviours (sponge-like media, aquifer, turbulence, etc.) namely finance, viscoelasticity, control theory, electrical networks, groundwater flow and geo-hydrology, wave propagation, plasma physics and fusion, rheology, chaotic processes, fluid mechanics and biological activities [1–6]. For more applications of fractal-fractional differential equations, see [7–10] and for recent results on fractional differential equations and their applications, see [11–13]. To explore more results on implicit fractional differential equations and their applications, see [14–17]. There are many results relating to implicit fractional differential equations in literature involving Caputo fractional derivatives both for initial value problems (IVP) and boundary value problems (BVP) [15,18–21].

In 2015, Benchohra and Souid in [18] studied the existence of integrable solutions for IVP for some given implicit fractional order functional differential equations with infinite delay

\[ \begin{cases} \mathcal{C}D^\mu \phi(\tau) = \sigma(\tau, \phi(\tau), \mathcal{C}D^\mu \phi(\tau)), \tau \in I = [0, b], \\ \phi(\tau) = \varphi(\tau), \quad \tau \in (-\infty, 0], \end{cases} \]

where \( \mathcal{C}D^\mu \) is the Caputo fractional differential operator, \( \sigma : I \times \mathbb{B} \times \mathbb{B} \to \mathbb{R} \) is a given function and \( \mathbb{B} \) is a phase space with its element \( \phi(\tau + \varepsilon) = \phi(\tau + \varepsilon), \quad \varepsilon \in [-\infty, 0] \).

In 2016, Kucche et al., in [20], considered the following equation:

\[ \begin{cases} \mathcal{C}D^\mu \phi(\tau) = \sigma(\tau, \phi(\tau), \mathcal{C}D^\mu \phi(\tau)), \quad \phi(0) = \phi_0 \in \mathbb{R}, \tau \in [0, T], \end{cases} \]
where \( ^CD^\mu(0 < \mu < 1) \) stands for the Caputo fractional derivative and \( \sigma : [0, T] \times \mathbb{R} \times \mathbb{R} \to \mathbb{R} \) is a known continuous function fulfilling some conditions. The authors investigated the well-posedness, interval of existence, and continuous dependence on the initial condition of solutions to Equation (1). Recently, in 2021, Shabbir et al. [17] worked on an implicit boundary value problem (BVP) involving an Atangana–Baleanu–Caputo (ABC) derivative of the form

\[
\begin{aligned}
\begin{cases}
\frac{\partial}{\partial \tau} \phi(t) = \sigma(\tau, \phi(\tau)), & 0 < \mu < 2, \tau \in I = [0, b], \\
\phi(0) = \phi_0, \phi(b) = \phi_1,
\end{cases}
\end{aligned}
\]

where \( \frac{\partial}{\partial \tau} \phi(t) \) denotes the ABC derivative of order \( 1 < \mu < 2 \) and \( \sigma : I \times \mathbb{R} \to \mathbb{R} \) is a continuous function. Here, the authors established the existence of solution, uniqueness of solution and stability of solution to the class of implicit BVPs (2) with an ABC type derivative and integral.

Motivated by some applications of the implicit fractal–fractional differential equation in modeling complex phenomena and systems in porous media with memory, and the result in [17], where the authors used the ABC derivative operator to study Equation (2); therefore, we generalize (2) for a class of fractal–fractional derivative operator known as the Mittag–Leffler law type kernel \( \psi \), 1 < \mu < 2, \tau \in I \) is a continuous function. Here, the authors investigated the well-posedness, exponential growth bound, and long-time behaviour of a solution to the class of implicit BVPs (2) with an ABC type derivative and integral.

The organization of the paper is as follows. In Section 2, we present the preliminaries; and in Section 3, we give the statements and proofs of the main results of the paper. Section 5 contains a brief summary of the paper.

2. Preliminaries

In this section, one gives some concepts that will be useful for the main result.

\[
\begin{aligned}
\begin{cases}
\int_a^b \frac{\partial}{\partial \tau} \phi(t) = \sigma(\tau, \phi(\tau)), & 0 < \mu < 2, \tau \in I = [0, b], \\
\phi(0) = \phi_0, \phi(b) = \phi_1,
\end{cases}
\end{aligned}
\]

Next, we define the norm of the solution \( \psi \) by

\[
\|\psi\| := \sup_{a \leq t \leq T} |\psi(t)|.
\]

The organization of the paper is as follows. In Section 2, we present the preliminaries; and in Section 3, we give the statements and proofs of the main results of the paper. Section 5 contains a brief summary of the paper.
Definition 2 ([22]). Suppose \( \phi : (a, b) \to \mathbb{R} \) is differentiable and \( 0 < \mu < 1 \). Then the ABC fractional derivative for function \( \phi \) of order \( \mu \) is defined as

\[
\left[ a \right]^{ABC} D_t^\mu \phi(t) = \frac{\mathbb{M}(\mu)}{1 - \mu} \int_a^t \phi'(\tau) E_\mu \left[ -\frac{\mu}{1 - \mu} (t - \tau)^\mu \right] d\tau,
\]

with \( \mathbb{M}(\mu) > 0 \) a normalization function satisfying \( \mathbb{M}(0) = \mathbb{M}(1) = 1 \).

Definition 3 ([1–3]). Let \( \phi : (a, b) \to \mathbb{R} \) be a differentiable function, suppose \( \phi \) is fractal differentiable in \( (a, b) \) with order \( \nu > 0 \). Then, the fractal–fractional derivative of \( \phi(t) \) of order \( \mu \) in Caputo sense with the Mittag–Leffler type kernel is defined by

\[
\left[ a \right]^{FFM} D_t^{\mu, \nu} \phi(t) = \frac{\Delta B(\mu)}{1 - \mu} \int_a^t \frac{d^\nu \phi'(\tau)}{d\tau^\nu} E_\mu \left[ -\frac{\mu}{1 - \mu} (t - \tau)^\mu \right] d\tau,
\]

with \( 0 < \mu, \nu \leq 1 \) and \( \Delta B(\mu) = 1 - \mu + \frac{\mu}{1 - \mu} \). The generalized form is given by

\[
\left[ a \right]^{FFM} D_t^{\mu, \nu, \theta} \phi(t) = \frac{\Delta B(\mu)}{1 - \mu} \int_a^t \frac{d^\nu \phi'(\tau)}{d\tau^\nu} E_\mu \left[ -\frac{\mu}{1 - \mu} (t - \tau)^\mu \right] d\tau, \quad 0 < \mu, \nu, \theta \leq 1,
\]

where

\[
\frac{d^\nu \phi'(\tau)}{d\tau^\nu} = \lim_{\tau \to t} \frac{\Phi^\nu(t) - \Phi^\nu(\tau)}{t^\nu - \tau^\nu}.
\]

Remark 1. When \( \nu = 1 \) in Equation (6), one obtains Equation (5).

Definition 4 ([1–3]). Let \( \phi : (a, b) \to \mathbb{R} \) be a continuous function. Then, the fractal–fractional integral of \( \phi \) with order \( \mu \) possessing Mittag–Leffler type kernel is defined as

\[
\left[ a \right]^{FFM} I_t^{\mu, \nu} \phi(t) = \frac{\mu^\nu}{\Delta B(\mu) \Gamma(\mu)} \int_a^t \tau^{\nu - 1} (t - \tau)^{\mu - 1} \phi(\tau) d\tau + \frac{(1 - \mu)\Gamma(\nu - 1)}{\Delta B(\mu)} \phi(t).
\]

Definition 5 ([23]). One defines the incomplete beta function by

\[
B_\tau(\mu, \nu) = \int_0^\tau t^{\mu - 1} (1 - t)^{\nu - 1} dt, \quad \tau \in [0, 1].
\]

It also has a representation in terms of a hypergeometric function given by

\[
B_\tau(\mu, \nu) = \frac{\tau^\nu}{\mu} 2F_1(\mu, 1 - \nu; \mu + 1; \tau).
\]

Definition 6 ([24]). The regularized incomplete beta function is given by

\[
I_\tau(\mu, \nu) = \frac{B_\tau(\mu, \nu)}{B(\mu, \nu)} = \frac{1}{B(\mu, \nu)} \int_0^\tau \tau^{\mu - 1} (1 - \tau)^{\nu - 1} d\tau,
\]

satisfying the following properties:

1. \( I_1(\mu, \nu) = I_1(\mu + 1, v - 1) + \frac{\nu^\nu (1 - \tau)^{\nu - 1}}{\mu B(\mu, \nu)}; \)
2. \( I_1(\mu, \nu) = I_1(\mu + 1, v + 1) - \frac{\nu^\nu (1 - \tau)^{\nu - 1}}{\nu B(\nu, \mu)}; \)
3. \( I_1(\mu, \nu) = I_1(\mu + 1, v) + \frac{\nu^\nu (1 - \tau)^{\nu - 1}}{\nu B(\nu, \mu)}; \)
4. \( I_1(\mu, \nu) = I_1(\mu, v + 1) - \frac{\nu^\nu (1 - \tau)^{\nu - 1}}{\nu B(\nu, \mu)}; \)
5. \( I_1(\mu, \nu) + I_1(\nu, \mu) = 1; \)
6. \( I_1(\mu, \nu) = 1 \text{ and } I_1(\mu, \nu) \in [0, 1]. \)
Lemma 1 ([25]). For all $\mu > 0$, $v \geq 1$, $0 \leq \tau \leq 1$, we have

1. $\frac{\tau^\mu(1-\tau)^{v-1}}{\mu B(\mu,v)} \leq I_1(\mu,v)$;
2. $I_1(\mu,v) \leq \frac{\tau^\mu}{\mu B(\mu,v)}$.

3. Main Results

This section starts with a Lipschitz condition on $\zeta(\cdot)$:

Condition 1. Let $0 < \text{Lip}_\zeta < \infty$. Given that $\xi_1, \xi_2, \xi_1', \xi_2' \in \mathbb{R}$, one has

$$|\zeta(\tau, \xi_1, \xi_1') - \zeta(\tau, \xi_2, \xi_2')| \leq \text{Lip}_\zeta(|\xi_1 - \xi_1'| + |\xi_1 - \xi_2'|),$$

with $\zeta(\tau, 0, 0) = 0$ and

$$|\zeta(\tau, \xi_1, \xi_1)| \leq \text{Lip}_\zeta(|\xi_1| + |\xi_1|). \quad (7)$$

Lemma 2. Let $0 < \text{Lip}_\zeta < 1$ and Condition 1 holds. Then $\mathbb{F}_t \mathbb{D}_t^{\mu,\nu}$ is a global Lipschitz continuous operator.

Proof. From Equations (3) and (7), we have

$$\left| \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(\tau) \right| = \left| \zeta\left(\tau, \psi(\tau), \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(\tau)\right) \right| \leq \text{Lip}_\zeta \left( |\psi(\tau)| + \left| \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(\tau) \right| \right),$$

and, therefore,

$$\left| \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(\tau) \right| \leq \frac{\text{Lip}_\zeta}{1 - \text{Lip}_\zeta} |\psi(\tau)|.$$

Furthermore, one obtains

$$\left| \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(\tau) - \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \phi(\tau) \right| \leq \frac{\text{Lip}_\zeta}{1 - \text{Lip}_\zeta} |\psi(\tau) - \phi(\tau)|.$$

3.1. Existence and Uniqueness Result

Here, we establish the well-posedness of solution to Equation (3). Now, define

$$A\psi(t) = \psi_a + \frac{(1 - \mu)v}{\AB(\mu)} t^{v-1} \zeta\left(t, \psi(t), \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(t)\right) + \frac{\mu v}{\AB(\mu) \Gamma(\mu)} \int_0^t \tau^{v-1} (t - \tau)^{\mu - 1} \zeta\left(\tau, \psi(\tau), \mathbb{F}_t \mathbb{D}_t^{\mu,\nu} \psi(\tau)\right) d\tau,$$

and obtain the following auxiliary results:

Lemma 3. Let $\psi$ be a solution satisfying Equation (4) and let Condition 1 be satisfied. Then, it follows that for all $\mu, v \in (0, 1]$ such that $\mu + v \geq 1$, we have

$$\|A\psi\| \leq c_1 + c_2 \frac{\text{Lip}_\zeta}{1 - \text{Lip}_\zeta} \|\psi\|,$$

where $c_1$ and $c_2 := \frac{\mu v}{\AB(\mu)} \left(1 - \mu\right) a^{v-1} + \frac{\mu v}{\Gamma(\mu + v)} T^{\mu + v - 1}$ are positive constants with $|\psi_a| \leq c_1$. 
\textbf{Proof.} By taking absolute value on the operator $\mathcal{A}$, we have

$$|\mathcal{A}\psi(t)| \leq |\psi_a| + \frac{1 - \mu}{\mathcal{A}B(\mu)} \|D^\mu\psi(t)\|_1 \left|\int_a^t \tau^{\mu-1} (t - \tau)^{\mu-1} \left| \left( I, \psi(t) \right)_{\mathcal{A}} FFM D_{\mathcal{A}}^{\mu,\nu} \psi(\tau) \right| d\tau\right|$$

$$\leq |\psi_a| + \frac{1 - \mu}{\mathcal{A}B(\mu)} \|D^\mu\psi(t)\|_1 \left|\int_a^t \tau^{\mu-1} (t - \tau)^{\mu-1} \left| \left( I, \psi(t) \right)_{\mathcal{A}} FFM D_{\mathcal{A}}^{\mu,\nu} \psi(\tau) \right| d\tau\right|.$$

Applying Condition 1 and $|\varphi_a| \leq c_1$, to obtain

$$|\mathcal{A}\psi(t)| \leq c_1 + \frac{1 - \mu}{\mathcal{A}B(\mu)} \|D^\mu\psi(t)\|_1 \left|\int_a^t \tau^{\mu-1} (t - \tau)^{\mu-1} \left| \left( I, \psi(t) \right)_{\mathcal{A}} FFM D_{\mathcal{A}}^{\mu,\nu} \psi(\tau) \right| d\tau\right|$$

From Lemma 2, we arrive at

$$|\mathcal{A}\psi(t)| \leq c_1 + \frac{1 - \mu}{\mathcal{A}B(\mu)} \|D^\mu\psi(t)\|_1 \left|\int_a^t \tau^{\mu-1} (t - \tau)^{\mu-1} \left| \left( I, \psi(t) \right)_{\mathcal{A}} FFM D_{\mathcal{A}}^{\mu,\nu} \psi(\tau) \right| d\tau\right|$$

Evaluating the integral above, we have

$$|\mathcal{A}\psi(t)| \leq c_1 + \frac{1 - \mu}{\mathcal{A}B(\mu)} \|D^\mu\psi(t)\|_1 \left|\int_a^t \tau^{\mu-1} (t - \tau)^{\mu-1} \left| \left( I, \psi(t) \right)_{\mathcal{A}} FFM D_{\mathcal{A}}^{\mu,\nu} \psi(\tau) \right| d\tau\right|$$

$$\leq c_1 + \frac{1 - \mu}{\mathcal{A}B(\mu)} \|D^\mu\psi(t)\|_1 \left|\int_a^t \tau^{\mu-1} (t - \tau)^{\mu-1} \left| \left( I, \psi(t) \right)_{\mathcal{A}} FFM D_{\mathcal{A}}^{\mu,\nu} \psi(\tau) \right| d\tau\right|.$$ (8)
since \( 1 - I_{\phi}^1(\mu, v) \in [0,1] \). We observe that \( a^{v-1} > t^{v-1} \) since \( v - 1 < 0 \). Thus, taking supremum over \( t \in [a, T] \) in (8) and recalling that \( \mu + v \geq 1 \), we obtain

\[
\|A\phi\| \leq c_1 + \frac{(1-\mu)v}{\lambda(\mu)} \|\psi\| + \frac{\mu v}{\lambda(\mu)} \|\psi\| + \frac{\mu v}{\lambda(\mu)} \|\psi\| \|B(\mu, v)T^{\mu+v-1}\|
\]

\[
eq c_1 + \frac{v}{\lambda(\mu)} \frac{\lambda_1}{\lambda_1 - \lambda_5} \left[ (1 - \mu)a^{v-1} + \frac{\mu(\mu)T^{\mu+v-1}}{\Gamma(\mu + v)} \right] \|\psi\|,
\]

and the proof is complete. \( \square \)

**Lemma 4.** Suppose \( \psi \) and \( \phi \) are solutions satisfying Equation (4) and let Condition 1 be satisfied. Then, if it follows that for all \( \mu, v \in (0, 1) \) such that \( \mu + v \geq 1 \), we have

\[
\|A\phi - A\psi\| \leq c_2 \|\psi - \phi\|.
\]

**Proof.** The proof is skipped since it follows similar steps as the proof of Lemma 3. \( \square \)

Next, we state the existence and uniqueness theorem for Equation (3).

**Theorem 1.** Let \( \alpha + \beta \geq 1 \) and suppose Condition 1 is satisfied. Let \( c_2 > 0 \), such that \( c_2 < \frac{1 - \text{Lip}_A}{\text{Lip}_c} \), where \( c_2 := \frac{v}{\lambda(\mu)} \left( (1 - \mu)a^{v-1} + \frac{\mu(\mu)T^{\mu+v-1}}{\Gamma(\mu + v)} \right) \). Then, there exists a unique solution to Equation (3).

**Proof.** We proceed by applying the Banach fixed point theorem. Let \( \psi(t) = A\psi(t) \), then using Lemma 3, we have

\[
\|\psi\| = \|A\psi\| \leq c_1 + c_2 \frac{\text{Lip}_c}{1 - \text{Lip}_c} \|\psi\|.
\]

Collecting similar terms, we have \( \|\psi\| \left( 1 - c_2 \frac{\text{Lip}_c}{1 - \text{Lip}_c} \right) \leq c_1 \). This shows that \( \|\psi\| < \infty \) since \( c_2 < \frac{1 - \text{Lip}_A}{\text{Lip}_c} \). Furthermore, if \( \psi \neq \phi \) are solutions to Equation (3), then from Lemma 4, we have

\[
\|\psi - \phi\| = \|A\psi - A\phi\| \leq c_2 \frac{\text{Lip}_c}{1 - \text{Lip}_c} \|\psi - \phi\|.
\]

It follows that \( \|\psi - \phi\| \left( 1 - c_2 \frac{\text{Lip}_c}{1 - \text{Lip}_c} \right) \leq 0 \). Since \( 1 - c_2 \frac{\text{Lip}_c}{1 - \text{Lip}_c} > 0 \), that is, \( c_2 < \frac{1 - \text{Lip}_A}{\text{Lip}_c} \), then \( \|\psi - \phi\| < 0 \). This is a contradiction and, therefore, \( \|\psi - \phi\| = 0 \). \( \square \)

**Exponential Growth**

We present an inequality needed in proving the upper growth bound:

**Theorem 2 ([26]).** Given that \( f, g, h : I \rightarrow \mathbb{R}^+ \) are continuous functions. If \( \phi : I \rightarrow \mathbb{R}^+ \) is continuous and

\[
\phi(t) \leq f(t) + \int_{t_0}^{t}(t - \tau)^{v-1}g(\tau)\phi(\tau)d\tau + \int_{t_0}^{t}(t - \tau)^{v-1}h(\tau)\phi^{\gamma}(\tau)d\tau, \quad t \in I,
\]

with constants \( v > 0 \) and \( 0 < \gamma < 1 \), then the following statements are true.
(i) Given that \( v > \frac{1}{2} \). It follows that
\[
\phi(t) \leq \left[ A_1^{1-\gamma}(t) \right. \\
+ (1 - \gamma)K_1 \int_0^t \exp \left( (\gamma - 1)K_1 \int_0^s g^2(\tau)d\tau \right) h^2(s)R_1(s)ds \left. \right] \frac{1}{s^{1-\gamma}} \\
\times \exp \left( t + \frac{K_1}{2} \int_0^t g^2(s)ds \right), \ t \in I,
\]
where \( A_1(t) = \max_{l_0 \leq l \leq t} \{ 2e^{-2s}f^2(s) \} \), \( K_1 = \frac{2(2v-1)}{4v-1} \), and \( R_1(t) = e^{2(v-1)t} \).

(ii) Given that \( \rho \in (0, \frac{1}{2}) \), \( q = \frac{1+v}{4} \) and \( p = 1 + v \). Then
\[
\phi(t) \leq \left[ A_2^{1-\gamma}(t) \right. \\
+ (1 - \gamma)K_2 \int_0^t \exp \left( (\gamma - 1)K_2 \int_0^s g^2(\tau)d\tau \right) h^2(s)R_2(s)ds \left. \right] \frac{1}{s^{1-\gamma}} \\
\times \exp \left( t + \frac{K_2}{q} \int_0^t g^2(s)ds \right), \ t \in I,
\]
where \( A_2(t) = \max_{l_0 \leq l \leq t} \{ 2e^{-1}e^{-se}f^2(s) \} \), \( K_2 = 2^{2e-2} \left( \frac{1}{p} \right) 2^{\frac{2v-1}{4v-1} - \frac{1}{v}} \), and \( R_2(t) = e^{2(v-1)t} \).

**Theorem 3.** Given that \( \psi \) satisfies Equation (4) and the initial function \( \psi_0 \) is bounded above. Suppose Condition 1 is satisfied, then it follows that for all \( \mu \in (\frac{1}{2}, 1] \) and \( v \in [0, 1] \), one gets
\[
|\psi(t)| \leq c_5 \exp \left( t - c_6 t^{v-1} \right), \ t \in [a, T],
\]
where \( c_5 = (2e^{-2k_c}c_3^{1/2}) \exp \left( \frac{d^{v-1}}{1-v} \right) \) and \( c_6 = (c_4 \text{Lip}_c)^2 \frac{2^{\frac{2v-1}{4v-1}} - \frac{1}{v}}{2^{\frac{2v-1}{4v-1}}} \) are some positive numbers.

**Proof.** Following from the line of proof of Theorem 3, one obtains
\[
|\psi(t)| \leq c_1 + \frac{(1 - \mu) \nu}{A_\mathcal{B}(\mu) \Gamma(\mu)} \int_0^t \frac{\text{Lip}_c}{1 - \text{Lip}_c} |\psi(t)| \\
+ \frac{\mu \nu}{A_\mathcal{B}(\mu) \Gamma(\mu)} \int_a^t \frac{\text{Lip}_c}{1 - \text{Lip}_c} |\psi(t)| d\tau \\
\leq c_1 + \frac{(1 - \mu) \nu}{A_\mathcal{B}(\mu) \Gamma(\mu)} \int_a^t \frac{\text{Lip}_c}{1 - \text{Lip}_c} \int_0^t \tau^{v-1}(t - \tau)^{\mu-1} |\psi(\tau)| d\tau \\
+ \frac{\mu \nu}{A_\mathcal{B}(\mu) \Gamma(\mu)} \int_a^t \frac{\text{Lip}_c}{1 - \text{Lip}_c} \int_0^t \tau^{v-1}(t - \tau)^{\mu-1} Y(\tau) d\tau. \tag{9}
\]
Let \( Y(t) := |\psi(t)| \) and \( \tilde{c}_2 := 1 - \frac{\text{Lip}_c}{\text{Lip}_c} \frac{d^{v-1}(1-\mu)\nu}{A_\mathcal{B}(\mu)} > 0 \). Thus, for \( \text{Lip}_c < \frac{A_\mathcal{B}(\mu) \nu}{A_\mathcal{B}(\mu) \nu (1-\mu) d^{v-1}} \), it follows from (9) that
\[
\tilde{c}_2 Y(t) \leq c_1 + \frac{\mu \nu}{A_\mathcal{B}(\mu) \Gamma(\mu)} \int_a^t \tau^{v-1}(t - \tau)^{\mu-1} Y(\tau) d\tau.
\]
Now, dividing by \( \tilde{c}_2 = \frac{\lambda B(\mu) - \text{Lip}_c}{(1 - \text{Lip}_c) A(\mu)} \), to get

\[
Y(t) \leq c_3 + c_4 \text{Lip}_c \int_0^t \tau^{v-1} (t - \tau)^{\mu-1} Y(\tau) d\tau,
\]

with \( c_3 := \frac{1}{\tilde{c}_2^2} \) and \( c_4 := \frac{\mu}{\Gamma(\mu)} (A(\mu) - \text{Lip}_c A(\mu + (1 - \mu) \alpha^\nu - 1)) \). Next, using Theorem 2 (i) with \( h(\tau) = 0, g(\tau) = c_4 \text{Lip}_c \tau^{v-1} \) and \( v = \mu > \frac{1}{2} \), to obtain

\[
Y(t) \leq A_1^{1/2}(t) \exp \left( t + \frac{K_1}{2} (c_4 \text{Lip}_c)^2 \int_0^t \tau^{2v-2} d\tau \right) = A_1^{1/2}(t) \exp \left( t + (c_4 \text{Lip}_c)^2 \frac{\Gamma(2\mu - 1)}{4\mu - 1} \frac{t^{\nu-1} - \frac{1}{4\mu - 1}}{v - 1} \right) \leq (2e^{-2e^2c_3^2})^{1/2} \exp \left( t - (c_4 \text{Lip}_c)^2 \frac{\Gamma(2\mu - 1)}{4\mu - 1} \frac{t^{\nu-1}}{1 - v} \right),
\]

where \( K_1 = \frac{2\Gamma(2\mu - 1)}{4\mu - 1} \) and \( A_1(t) = \max_{a \leq t \leq t} (2e^{-2e^2c_3^2}) = 2e^{-2e^2c_3^2} \), since \( e^{-2t} \) is decreasing. \( \square \)

3.2. Asymptotic Property of the Solution

Here, we show the long term (limiting) property of our solution. The corollary indicates that the rate of energy growth of the solution is finite when time becomes large.

Corollary 1. Under the hypotheses of Theorem 3 and for all \( 0 < \beta < 1 \), we have

\[
\limsup_{t \to \infty} \frac{\log |\psi(t)|}{t} \leq 1.
\]

Proof. We obtained from Theorem 3 that

\[
|\psi(t)| \leq c_5 \exp \left( t - c_6 t^{\nu-1} \right), t \in [a, T].
\]

If we take log of both sides of the above equation, it will yield

\[
\log |\psi(t)| \leq \log(c_5) + t - c_6 t^{\nu-1}.
\]

Next, divide through by \( t \) to obtain

\[
\frac{\log |\psi(t)|}{t} \leq \frac{\log(c_5)}{t} + 1 - c_6 \frac{t^{\nu-1}}{t} = \frac{\log(c_5)}{t} + 1 - \frac{c_6}{t^{\nu-1}}.
\]

Since \( 0 < \nu < 1 \), it follows that \( 2 - \nu \geq 1 \). Now, take limit supremum over \( t \) in both sides to get

\[
\limsup_{t \to \infty} \frac{\log |\psi(t)|}{t} \leq \limsup_{t \to \infty} \frac{\log(c_5)}{t} + 1 - \limsup_{t \to \infty} \frac{c_6}{t^{2-\nu}} = 1.
\]

\( \square \)

4. Examples

Now, we give examples to illustrate the result in Theorem 3. The following are some plots (graphs) for the upper bound growth of our energy solution \( |\psi(t)| \leq \exp \left( t - \frac{1}{1-v} t^{\nu-1} \right), t \in [a, T] \). For convenience, we set \( c_5 = 1 \) and choose \( \mu \in \left( \frac{1}{2}, 1 \right) \), such that \( c_6 = \frac{1}{1-v} \) with \( c_4 = 1, \text{Lip}_c = 1 \). We plotted graphs of the growth bound for \( v = \frac{1}{3}, \frac{1}{5}, \frac{1}{7}, \frac{1}{9}, \frac{2}{5}, \frac{2}{7}, \frac{2}{9}, \frac{3}{9}, \frac{4}{9}, \frac{5}{9}, \frac{6}{9}, \frac{7}{9}, \frac{8}{9}, \frac{9}{9} \), and for various time intervals. It is observed that as closer the pa-
rameter $\nu$ is to zero, the faster the rate of growth is to the bound. However, as time grows large, the growth rate is at most at $t = 600$ irrespective of the values of $\nu$, as shown in the Figure 1 below.

Some graphs of the upper growth bound: $\exp(t - \frac{1}{1-\beta} t^{\beta-1})$ for $\beta = \frac{1}{10}, \frac{1}{5}, \frac{1}{4}, \frac{1}{3}, \frac{2}{10}, \frac{8}{10}, \frac{9}{10}$, and for various time intervals.

Figure 1. Graphical illustration of the growth bounds.

5. Conclusions

Fractional order derivatives are used to represent memory formalism in modeling phenomena or processes in porous media in order to diminish the size of the pores and the permeability of the porous matrix [27]. Hence, implicit fractal–fractional differential equations are very important because they model many technical processes and systems in porous environment exhibiting long time memory property. In this paper, we estimated the higher growth bound of our solution and it is shown that the solution exhibits an exponential growth in $t$ at a specific rate. Furthermore, the result shows a long time behaviour of the mild solution. Banach fixed point theorem was applied to prove the well-posedness of mild solution to the class of implicit time–fractal–fractional differential
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