BEINIT: Avoiding Barren Plateaus in Variational Quantum Algorithms

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Abstract—Barren plateaus are a notorious problem in the optimization of variational quantum algorithms and pose a critical obstacle in the quest for more efficient quantum machine learning algorithms. Many potential reasons for barren plateaus have been identified but few solutions have been proposed to avoid them in practice. Existing solutions are mainly focused on the initialization of unitary gate parameters without taking into account the changes induced by input data. In this paper, we propose an alternative strategy which initializes the parameters of a unitary gate by drawing from a beta distribution. The hyperparameters of the beta distribution are estimated from the data. To further prevent barren plateau during training we add a novel perturbation at every gradient descent step. Taking these ideas together, we empirically show that our proposed framework significantly reduces the possibility of a complex quantum neural network getting stuck in a barren plateau.

Reproducibility: The source code and data are available at https://github.com/aicaffeinelife/BEINIT

I. INTRODUCTION

The world of machine learning has seen a tectonic paradigm shift in the last decade. From handcrafted features that were manually fed to classical optimization algorithms like support vector machines we have moved on to the deep neural networks (DNN) that can simultaneously perform the task of feature engineering and recognition without any explicit intervention from the developer. The ever increasing performance gain in terms of accuracy on really large datasets has led to the quest for more faster and accurate algorithms that can scale to the amount of data being produced today.

The ongoing development of Noisy Intermediate-Scale Quantum (NISQ) computers has led to considerable excitement about the potential quantum advantage that can be obtained in several optimization problems that are used in almost all fields of science today. Various hybrid quantum-classical algorithms [1]–[4] offer possible ways to utilize NISQ computers. Variational Quantum Algorithms (VQAs) offer a potential way of leveraging quantum computers alongside classical computers in a hybrid fashion. When it comes to deployment of VQAs on NISQ devices, the limitations of existing quantum computers (i.e., noisy gates and limited circuit depth) [2], [5] restrict the overall potential of these algorithms. Moreover, VQAs themselves are not without issues. In this paper, we make a step towards addressing a well known issue that frequently arises in optimization approaches involving VQAs.

The main component of a VQA is a variational quantum circuit (VQC) that consists of a finite sequence of parameterized unitary gates $U(\bar{\theta})$ and a measurement produced by a quantum observable $\hat{O}$. The parameters of the unitary gates are real-valued. More formally, we can express $U(\bar{\theta})$ as:

$$U(\bar{\theta}) = \prod_{i=1}^{L} U(\theta_1)U(\theta_2) \ldots U(\theta_L)$$ (1)

An ansatz formed of a sequence of unitary gates as in Equation 1 is trained to solve an optimization problem of the form

$$\theta^* = \arg \min_{\theta} C(\theta),$$

where $C(\theta)$ is a cost function that quantifies the deviation of a VQC at a given timestep from the “true” trajectory of the problem. Given $\hat{O}$ and $\theta$ a general cost function for VQAs can be defined as:

$$C(\theta) = f(|\langle 0 | U^\dagger(\bar{\theta}) \hat{O} U(\bar{\theta}) | 0 \rangle|)$$ (2)

Where the choice of $f$ depends on the specific application of the VQA [2]. The parameters of a VQC are initialized from a random distribution. Over the course of training, the output of the VQC is measured with respect to the observable $\hat{O}$ and the gradients of every unitary gate are estimated using the parameter shift rule [6], [7]. The set of parameters at a given time step and the gradients of the unitary gates are then passed to a classical computer which updates these parameters according to a gradient descent rule. This update is usually performed by an optimizer and it’s choice is usually considered to be a hyperparameter (i.e., parameters which are manually selected by the developer). Some popular optimizers include Adam [8], AdaGrad [9] etc.

Ideally, for an optimization problem, the set $\theta^*$ should correspond to the best set of parameters that minimizes the cost function. However, it has been observed that when the VQC is complex either in number of qubits used to represent the input or the number of layers of unitaries, the optimization halts at a suboptimal set of parameters that do not correspond to a minima in the optimization surface. This situation occurs when the circuit gets stuck on a plateau from where there are no good descent directions and is commonly referred to as being stuck in a “barren-plateau”. The unitary designs induced due to random parametrization of ansatz [10], optimization w.r.t a global cost function [11], entanglement between visible and hidden layers of a Quantum Neural Network (QNN) [12] and random entanglement [13] have been studied as potential causes for the barren plateau problem. Attempts
to mitigate the barren plateau problem comprise of advanced initialization schemes [14, 15], local objective functions [11] and Bayesian initialization schemes [16].

Our contribution The current body of work leaves, among the others, two questions unanswered. First, does the choice of the initialzation distribution have an effect on the appearance of barren plateaus in a random ansatz (barring the case where the parameters are initialized from a classical neural network)? Second, once a good initialization has been found, is there a way to influence the gradient descent such that the likelihood of barren plateaus is reduced? In this paper, we study these questions for the specific case of QNNs and empirically show that the choice of initialization distribution is an important and often neglected hyperparameter in VQAs. We further show that adding a derived perturbation in the parameter space during gradient descent can help a QNN from getting stuck in barren plateaus. We combine these insights into an algorithm which we call BEINIT. Although our results are derived for the specific case of QNN, we believe that these results are general enough for any VQA application.

II. BAREN PLATEAU PROBLEM IN QNNs

We earlier mentioned that barren plateaus can have several causes. In this work, we choose to focus on a generic case where we only vary the parameters of the unitary gates (as opposed to additionally randomizing the specific gate choice as in [10]) and optimize in presence of a global cost function i.e. for a given QNN we have one global observable. The global cost function assumption is not unreasonable since we wish to study the effect of initialization in cases where no additional steps to mitigate barren plateaus have already been taken (e.g. optimizing with local observables [11]).

Since we deal with QNNs, we briefly specify the general VQA framework discussed earlier. For any labeled dataset consisting of m samples; \(D = \{(x_i, y_i)\}_{i=1}^{m}\), a QNN requires access to a state encoding function \(K(\vec{z})\) which accepts \(\vec{z}^{(i)} \in \mathbb{R}^d\) and produces a quantum state. This encoding can be seen as an instance of kernel function which maps a real valued vector into the complex Hilbert space. The output state before measurement of a QNN can be described as \(|\psi\rangle = U(\vec{\theta})K(\vec{z})|0\rangle\).

The output of a classifier \(f(\vec{x}, \vec{\theta}, \vec{O})\) is an expectation over \(m\) samples in the given dataset \(f(\vec{x}, \vec{\theta}, \vec{O}) = \sum_{i=1}^{m} ||y_i - \langle 0|U^\dagger(\vec{\theta})K(\vec{z}^{(i)})\vec{O}U(\vec{\theta})K(\vec{z}^{(i)})|0\rangle||^2\).

This output is a scalar value which is compared using a distance measure with the true label and the error signal is propagated to each unitary circuit. An advantage of using such differentiation rule as in [6] is that the same unitary circuit can evaluate the output and its respective gradient. A circuit with unitary gates of dimension \(d\) corresponds to a unitary group \(U(d)\). For any measure \(dU\) on this group, the \(k^{th}\) moments can be expressed as in [17] \(M_k(dU) = \int_{U(d)} dU U_{ik,1j} \ldots U_{ik,jk} U_{1i',j'i'} \ldots U_{1k,k'}\), where \(ik, jk\) are the row and column indices of the \(k^{th}\) unitary matrix. If the measure is a left invariant measure like Haar-measure, i.e. \(dU = dU_H\) then, the first and second order moments are given as

\[
M_1(dU_H) = \frac{1}{d} \delta_{i1} \delta_{j1},
\]

and

\[
M_2(dU_H) = \frac{\delta_{i1} \delta_{j1} \delta_{j1} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2}}{d^3} - \frac{\delta_{i1} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2} \delta_{j2}}{d(d^2 - 1)},
\]

respectively. Any measure \(dU\) exhibits a unitary 2-design if and only if \(M_1(dU) = M_1(dU_H)\) and \(M_2(dU) = M_2(dU_H)\).

For a sequence of parameterized unitary gates expressed as \(U(\vec{\theta}) = \prod_{i=1}^{n} U(\theta_1) \ldots U(\theta_L)\) we can denote

\[
U_- = \prod_{i=1}^{k} U(\theta_1) \ldots U(\theta_{k-1}) \quad \text{and} \quad U_+ = \prod_{i=1}^{k} U(\theta_k) \ldots U(\theta_L)
\]

for the \(k^{th}\) unitary \(U_k\).

In [10] the authors show that if any \(U_-\) or \(U_+\) exhibits a unitary 2-design then the variance of the classifier with respect to \(\theta_k\), \(\text{Var}[\partial_{\theta_k} f]\) can be given by:

\[
\text{Var}[\partial_{\theta_k} f] \propto \left( \frac{\text{Tr}(U^2)}{2^n} - \frac{\text{Tr}(U)^2}{2^{4n}} \right)
\]

In Equation (3), \(U \in U(d)\) and \(n\) is the number of qubits. The relation assumes that both \(U_-\) and \(U_+\) exhibit a 2-design and shows that with increasing number of qubits, the variance of the gradient will approach zero and consequently the QNN will be stuck in barren plateau with all conditions kept the same (other cases also exhibit a similar kind of decay).

III. AN EXPERIMENT WITH INITIALIZATION

We noted that related works are relying on the common assumption that the parameters for \(l^{th}\) layer unitary \(U_l(\theta_l)\) are drawn from a uniform distribution (e.g. from \([-\pi, \pi]\)). In order to answer our first question i.e. the effect of initialization on variance of gradient with different probability distributions, we reproduced the experimental setup of [10] and ran it with parameters derived from the Uniform (\(\theta_l \sim \text{Unif}(0, 2\pi)\)), Beta (\(\theta_l \sim \text{Beta}(1, 2\pi)\)) and a normal distributions (\(\theta_l \sim N(0, 2\pi)\)). The results of our experiment are shown in Figure 1b.

Note that the blue line corresponding to the uniform distribution closely matches the results obtained in [10]. Interestingly, when the parameters of the quantum circuit are initialized using the beta distribution (orange line), the reduction in variance of the gradient is significantly less than the other two cases. This result may seem counter-intuitive at first, but if we examine the Beta distribution more closely we find that its domain is \([0, 1]\) which implies that it can be used to model a distribution over probabilities. Moreover, its shape parameters \(\alpha, \beta\) (which are inputs to the gamma function \(\Gamma(z) = \int_0^\infty x^{z-1}e^{-x}dx\), see Eq. 4 and 5) can be tuned and used to guide the relative weight over the input probabilities.
This behavior is in contrast with the uniform distribution which places equal weight over the input interval.

Based on our experimental observations above we advocate for the following conjecture.

**Conjecture 1.** The choice of the sampling distribution over a unitary group $U(d)$ has a direct impact on the likelihood of the resulting QVC exhibiting a unitary-2 design. The more non-mean collapsing behavior a distribution exhibits, the lesser the likelihood of the occurrence of a barren plateau.

By a “mean-collapsing” distribution we refer to a distribution for which its first and second order moments are unbiased MLE estimators (e.g., normal distribution). Conversely, a non-mean collapsing distribution has biased estimators in either first or second order moments.

IV. EMPIRICAL BAYES

The effectiveness of Beta distribution raises a question - is it possible to find the shape parameters in a principled manner? We analyze this question using insights from the Bayesian machine learning.

For any given dataset $D$ and a model with parameters $\theta$, the goal of Bayesian machine learning is to determine the posterior distribution $\pi(\theta|D)$. A prior distribution over the parameters is given by $\pi(\theta)$ and the likelihood function is given as $L(D; \theta) = \prod_{i=1}^{m} p(x_i|\theta)$ where $x_i \in D$. By the Bayes Law $\pi(\theta|D) \propto L(D; \theta) \pi(\theta)$.

In a purely Bayesian model, we can set the prior $p(\theta)$ distribution to be a known parametrized distribution. We can further assume a prior on the hyperparameters of the parametrized distribution as well. However, doing this in practice is intractable for large dimensions or data points. A technique proposed in [18], [19] alleviates this problem by estimating the hyperparameters from the data using Maximum Likelihood Estimation (MLE) approach. These initial hyperparameters are then used to draw samples from the prior distribution. We use this insight to impose a known prior distribution over the parameters of the unitary gates and estimate the shape parameters directly from the data.

For the Beta distribution the MLE estimation of $\alpha$ and $\beta$ can be found as:

$$\alpha_{MLE} \equiv m \Gamma'(\alpha + \beta) - m \frac{\Gamma'(\alpha)}{\Gamma(\alpha)} + \sum_{i=1}^{m} x_i = 0 \quad (4)$$

$$\beta_{MLE} \equiv m \frac{\Gamma'(\alpha + \beta)}{\Gamma(\alpha + \beta)} - m \frac{\Gamma'(\beta)}{\Gamma(\beta)} + \sum_{i=1}^{m} (1 - x_i) = 0 \quad (5)$$

**Algorithm 1** The BEINIT Algorithm

1: procedure BEINIT ($D, \eta, C(\theta)$, iters)
2: $\alpha, \beta \leftarrow$ EB_FIT($D$)
3: $D_{\text{train}}, D_{\text{test}} \leftarrow$ SPLIT($D$)
4: $\theta_1 \leftarrow$ INIT($\alpha, \beta$)
5: $\theta \leftarrow \theta_1$
6: for $i = 0 \rightarrow$ iters do
7: $\sigma_\theta^2 \leftarrow \text{VAR} (\nabla_\theta L)$
8: $\eta \leftarrow \frac{1}{\sqrt{2}} (\frac{1 + \beta}{\beta} + \sigma_\theta^2)$
9: $\theta' \leftarrow \theta + \mathcal{N}(0, \eta)$
10: $\theta^{i+1} \leftarrow \text{GRAD_DESCENT}(\theta', D_{\text{train}}, \eta, C(\theta))$
11: end for
12: end procedure

V. THE BEINIT ALGORITHM

We propose an algorithm termed BEINIT that is based on two key observations. First, drawing samples from a beta distribution for parametrizing the quantum circuit yields a much higher gradient variance for increasing number of qubits and second, it is possible to estimate the hyperparameters of the beta distribution from the data itself. The second observation is based on the empirical Bayes framework discussed earlier.

Algorithm 1 describes the overall procedure and is visually represented in Figure 1a. We take as input a labelled dataset
\[\mathcal{D} = \{(x_i, y_i)\}_{i=1}^{m},\] the learning rate \(\eta,\) a cost function \(C(\theta)\) and the number of iteration steps \(\text{iters}\) as a required input. In the first step, we estimate the \(\alpha\) and \(\beta\) parameters from all the available data using Equations (4) and (5). We then split this dataset into training and test subsets and initialize the circuit parameters from the discovered beta distribution parameters.

During the training, we keep track of the gradient variance produced during an iteration \(i;\) \(\sigma_i^2.\) It has been shown that adding a perturbation during gradient descent can help a learning algorithm escape saddle points [20] [21]. Inspired by a perturbation method proposed for deep learning [22], we introduce noise in the parameter space by considering a normal distribution parameterized by 0 mean and \(\sigma_n^2\) variance given by:

\[
\sigma_n^2 = \frac{\eta}{(1 + i)(\gamma + \sigma_i^2)}, \tag{6}
\]

where \(\eta\) and \(\gamma\) are fixed parameters. One can interpret \(\gamma\) as a constant additive bias that prevents a low noise perturbation (since \(\sigma_i^2\) progressively decreases as the output of the circuit gets closer to the true label). Empirically, the choice of \(\eta\) defines the amount of noise that will be generated per gradient descent step. We follow earlier works and choose \(\eta \in \{0.01, 0.3, 1.0\}\) and keep \(\gamma = 0.55.\) It must be noted that our method generates this perturbation in the parameter-space as opposed the gradient space which is proposed by earlier related work [22].

VI. NUMERICAL EXPERIMENTS

In this section, we demonstrate numerical results that lend experimental evidence for our claims in previous sections. The source code and data are available at https://github.com/iacaffeinelife/BEINIT

A. Data Driven Initialization with Bayes Distribution

In our exposition of the BEINIT algorithm we had claimed that a data driven initialization process worked best with the Beta distribution. To verify this claim we extracted a two class subset from the original Iris dataset [23]. We then normalized the data and estimated the parameters of three continuous distributions - Beta, Uniform and Normal using MLE estimation similar to Equations (4) and (5). The results of this experiment are shown in Figure 2.

The solid histogram is the normalized subset of the Iris dataset. The histogram contoured with red line is constructed by drawing an equivalent number of samples from the three continuous random distributions with their estimated parameters. We can see that for the case of uniform distribution, the estimated upper and lower bounds correspond to the minimum and maximum value of the data and the random distribution is evenly spread around these bounds. A similar observation can be made about the case of the normal distribution where the random distribution concentrates heavily around the mean of the underlying data. In contrast to these two cases, the Beta distribution closely conforms to the contours of the data. These results indicate that the Beta distribution is a much more data-sensitive distribution and can help provide a parameter initialization such that the parameters do not initially concentrate around the first or second order moments of the underlying data.

B. Experiments with Increasing Number of Qubits

We demonstrate the effectiveness of our algorithm on the Iris and Wine datasets [23]. To study the effect of the BEINIT algorithm on the gradient variance with increasing number of qubits, we compare against two baseline cases. In the first case, we initialize the QNN’s parameters using the uniform distribution as per [10] with the parameters estimated from the normalized data distribution. In the second case, we keep the same initialization as before, but introduce the perturbation as in Equation (6). Both datasets considered in this study are multi-class classification problems, but we convert them into a binary classification problem by taking only first two classes. The labels are then binarized to \(\pm 1.\) In the case of the Wine dataset, we reduce the dimensionality of the data by performing a principal component analysis (PCA) such that \(d = 2,\) where \(d\) indicates the dimensionality of a given data vector. We choose the number of qubits \(q \in \{4, 5, 6, 7, 8, 9, 10\}\). All our experiments are performed on the default quantum simulator provided by the Pennylane [24] library and we use the Nesterov momentum optimizer [25].
The variance of the gradient of the first parameter is shown with increasing qubit sizes for three different initialization scenarios. Initialization with a beta distribution and added perturbation show the least degradation in variance with increasing qubit size.

Fig. 4: Training costs for the Wine dataset for different initialization schemes with and without parameter perturbation.

Figure 3 shows an interpolated line computed from the raw gradient values for different cases. We can observe that for both datasets, initializing with a beta distribution proves to be beneficial in reducing the likelihood of a 2-design from forming. We can also see that in certain cases (e.g. Wine dataset), perturbing the gradient even with uniform distribution can have a beneficial effect on the variance of the gradient. It is interesting to note that in both datasets, a variational quantum circuit initialized with uniform distribution exhibits a similar decay in variance as with a random circuit of [10]. This strongly indicates that a VQCs initialized with a uniform distribution have a higher likelihood of exhibiting a unitary 2-design.

Figure 4 further shows the training costs corresponding to the different cases considered earlier for the Wine dataset. We can see that adding perturbation to the parameters during training allows for more variation in the cost function even though the ansatz converges to almost the same end cost. This variation can thus allow an ansatz to escape a barren plateau by providing a good descent direction during training. The behaviour of other datasets is similar.

The QNN architecture used in our studies is shown in Figure 5. The first stage in our architecture is a standard state encoding circuit that maps \( \vec{x} \in \mathbb{R}^d \mapsto |\psi(x)\rangle \in \mathbb{C}^q \). For each dimension \( j = 0 \ldots d; d \in \{2, 4\} \), we encode \( x_j \in \vec{x} \) as \( |\psi(x)\rangle^{(j)} = e^{-i\pi x_j \sigma_x} \) where \( |\psi(x)\rangle^{(j)} \) indicates the \( j \)-th component of the resulting input quantum state and the RHS is the X rotation gate. Since we wanted to study the effect of increasing number of qubits, we kept the data encoding process as simple as possible. The rest of our ansatz contains a cascade of parameterized rotation gates organized in \( L \) layers. Their outputs are entangled using CNOT gates. We apply a Pauli-Z operator as the basis for measuring the output state \( |\psi_o\rangle \).

For all experiments.
C. Experiments with Deeper Circuits

![Graphs showing variance of gradient with increasing number of layers for Iris and Wine datasets.]

We also performed experiments with QNNs that were several layers deep on the two datasets above. In this set of experiments, we kept the number of qubits to be constant at 4 and varied the number of layers $L \in \{2, 4, 6, 8, 10, 15, 20, 25, 30\}$. Besides these changes, we kept the experimental setup similar to the qubit experiments before.

The variance of gradient of the first parameters in different cases is shown in Figure 6 shows the result of our experiments. As expected, the uniform distribution with no perturbation shows a poor scaling with increasing number of layers. With parameter perturbation, both uniform and beta initialization show a more stable behavior. In the case of the Iris dataset, the BEINIT procedure results in a steadier variance as compared to the uniform initialization with perturbation. The performance of BEINIT is also evidenced in the case of the Wine dataset. The results of these experiments shed light on the efficacy of both perturbation and data driven beta initialization for deep QNN circuits.

VII. DISCUSSION

In this paper, two crucial questions regarding solving the barren plateau problem have been raised. We first explored different probability distributions and found that the choice of probability distribution for drawing the initial set of parameters has a direct influence on the variance of the gradient with complex circuits (i.e., circuits with a large number of qubits). Our experiments show that the Beta distribution is a highly data-sensitive distribution and initializing from it can lead to much lower degradation in gradient variance. This result is general and can be applied to initialization of parameters in different VQA algorithms. In the case of QNNs, the data guides the end result of optimization. We thus design a data-driven initialization process that incorporates the insights from data (e.g., determining the hyperparameters of the initializing distribution) to guide the overall initialization process. The observations from our experiments confirm that this data-driven initialization process can have a strong positive influence on the overall training process and reduce the chances of a barren plateaus. Our experimental insights lead us to develop a conjecture that explores the relationship between barren plateaus in optimization and the characteristics of the initializing distribution. To the best of our knowledge, this is the first work that considers such a relationship in the context of VQA.

Our second question opened a novel direction in the exploration of solutions to avoid the barren plateau problem. In the previous works, no attention has been paid to the effect of perturbation during the course of training and its subsequent role in improving the variance of the circuit. We devised the BEINIT procedure to combine our earlier insights along with a parameter perturbation scheme and showed that it is effective in improving the variance of the circuit with increasing number of qubits. Our experiments with the initialization and perturbation are shedding light on the barren plateau problem and suggest that a solution to overcome it may lie in finding a good data driven initialization and perturbation during the course of gradient descent. Another promising research direction is to investigate how the proposed technique affects another variational quantum approach, namely, the quantum approximate optimization algorithm (QAOA). Its acceleration at increasing circuit depth and number of qubits including techniques to deal with the barren plateau is a subject of major efforts [26], [27].

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REFERENCES

[1] E. Farhi, J. Goldstone, and S. Gutmann, “A quantum approximate optimization algorithm,” preprint at https://arxiv.org/abs/1411.0408 (2014).

[2] M. Cerezo, A. Arrasmith, R. Babbush, S. C. Benjamin, S. Endo, K. Fujii, J. R. McClean, K. Mitarai, X. Yuan, L. Cincio et al., “Variational quantum algorithms,” Nature Reviews Physics, vol. 3, no. 9, pp. 625–644, 2021.

[3] H. Ushijima-Mwesigwa, R. Shaydulin, C. F. Negre, S. M. Mniszewski, Y. Alexeev, and I. Safro, “Multilevel combinatorial optimization across quantum architectures,” ACM Transactions on Quantum Computing, vol. 2, no. 1, pp. 1–29, 2021.

[4] R. Shaydulin, H. Ushijima-Mwesigwa, C. F. Negre, I. Safro, S. M. Mniszewski, and Y. Alexeev, “A hybrid approach for solving optimization problems on small quantum computers,” Computer, vol. 52, no. 6, pp. 18–26, 2019.

[5] X. Liu, A. Angone, R. Shaydulin, I. Safro, Y. Alexeev, and L. Cincio, “Layer VQE: A variational approach for combinatorial optimization on noisy quantum computers,” IEEE Transactions on Quantum Engineering, vol. 3, pp. 1–20, 2022.

[6] M. Schuld, V. Bergholm, C. Gogolin, J. Izaac, and N. Killoran, “Evaluating analytic gradients on quantum hardware,” Phys. Rev. A, vol. 99, p. 032331, Mar 2019. [Online]. Available: https://link.aps.org/doi/10.1103/PhysRevA.99.032331

[7] K. Mitarai, M. Negoro, M. Kitagawa, and K. Fujii, “Quantum circuit learning,” Phys. Rev. A, vol. 98, no. 3, p. 032309, 2018. [Online]. Available: https://journals.aps.org/pra/abstract/10.1103/PhysRevA.98.032309

[8] D. P. Kingma and J. Ba, “Adam: A method for stochastic optimization,” arXiv preprint arXiv:1412.6980, 2014.

[9] J. Duchi, E. Hazan, and Y. Singer, “Adaptive subgradient methods for online learning and stochastic optimization.” Journal of machine learning research, vol. 12, no. 7, 2011.

[10] J. R. McClean, S. Boixo, V. N. Smelyanskiy, R. Babbush, and H. Neven, “Barren plateaus in quantum neural network training landscapes,” Nature communications, vol. 9, no. 1, p. 4812, 2018. [Online]. Available: https://www.nature.com/articles/s41467-018-07090-4

[11] M. Cerezo, A. Sone, T. Volkoff, L. Cincio, and P. J. Coles, “Cost-function-dependent barren plateaus in shallow quantum neural networks,” preprint at https://arxiv.org/abs/2001.00550 (2020).

[12] C. O. Marrero, M. Kieferová, and N. Wiebe, “Entanglement-induced barren plateaus,” PRX Quantum, vol. 2, no. 4, p. 040316, 2021.

[13] T. L. Pati, K. Najafi, X. Gao, and S. F. Yelin, “Entanglement devised barren plateau mitigation,” Physical Review Research, vol. 3, no. 3, p. 033090, 2021.

[14] E. Grant, L. Wossnig, M. Ostaszewski, and M. Benedetti, “An initialization strategy for addressing barren plateaus in parametrized quantum circuits,” Quantum, vol. 3, p. 214, 2019. [Online]. Available: https://quantum-journal.org/papers/q-2019-12-09-214

[15] Verdon, Guillaume and Broughton, Michael and McClean, Jarrod R and Sung, Kevin J and Babbush, Ryan and Jiang, Zhang and Neven, Hartmut and Mohseni, Masoud, “Learning to learn with quantum neural networks,” preprint at https://arxiv.org/abs/1411.4028 (2014).

[16] R. Shaydulin, I. Safro, and K. Larson, “Multistart methods for quantum approximate optimization,” in 2019 IEEE High Performance Extreme Computing Conference (HPEC). IEEE, 2019, pp. 1–8. [Online]. Available: https://ieeexplore.ieee.org/document/8916288/

[17] D. Dua and C. Graff, “UCI machine learning repository,” 2017. [Online]. Available: http://archive.ics.uci.edu/ml

[18] M. Cerezo, A. Arrasmith, R. Babbush, S. C. Benjamin, S. Endo, K. Fujii, J. R. McClean, K. Mitarai, X. Yuan, L. Cincio et al., “Variational quantum algorithms,” Nature Reviews Physics, vol. 3, no. 9, pp. 625–644, 2021.

[19] Y. E. Nesterov, “A method for solving the convex programming problem with convergence rate o (1/k^2),” in Dokl. akad. nauk Ssr, vol. 289, 1983, pp. 543–547.

[20] A. Neelakantan, L. Vilnis, Q. V. Le, I. Sutskever, L. Kaiser, K. Kurach, and J. Martin, “Adding gradient noise improves learning for very deep networks,” arXiv preprint arXiv:1511.06907, 2015.