A SECOND LOOK AT THE METAL LINE VARIABILITY OF G29-38

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ABSTRACT

The pulsating white dwarf G29-38 possesses a dust disk and metal lines attributed to the accretion of its disk material. Von Hippel and Thompson have reported variability in the equivalent width of G29-38’s Ca ii K line on a timescale of days. We use high-resolution optical spectroscopy of G29-38’s Ca ii K line to test this observation. Over 6 days spanning 2007 June to 2007 October we see no evidence for variability in the equivalent width of the Ca ii K line. We also sample the variability of the Ca ii K line over integrated timescales of ~100–500 s, where errors from incomplete coverage of pulsation modes are predicted to be ~8%–15%. We find that the scatter of the equivalent widths over this time period is consistent with measurement errors at the 7% level, slightly weaker than predicted but within the uncertainties of predictions. Weaker Ca and Mg lines observed show no significant variability on yearly timescales over 10 years based on our data and other high-resolution spectra. We conclude that further study is warranted to verify whether the accretion onto G29-38 is variable.

Subject headings: circumstellar matter — stars: individual (G29-38)

1. INTRODUCTION

G29-38 is a fascinating white dwarf. It is at a distance of 14 pc (van Altena et al. 2001), it pulsates (Kleinman 1998), it shows an atmosphere polluted with Ca, Mg, and Fe (Koester et al. 1997), and it possesses an infrared excess attributed to a dust disk within its tidal disruption radius (Zuckerman & Becklin 1987; Jura 2003; Reach et al. 2005). Since tidally disrupted planetesimals are a likely explanation for the dust disk and G29-38’s metal pollution, it is tempting to speculate on the presence of a planet or planets perturbing objects close to the white dwarf (Debes & Sigurdsson 2002). This perturbation, which could be constant, periodic, or stochastic, may cause variability in the amount of material that drizzles onto the atmosphere of G29-38 from the disk.

Several factors complicate the search for accretion variability in G29-38. G29-38 is a ZZ Ceti pulsator at the cool end of the DA variability strip, with peak-to-peak continuum changes of a few percent (Winget et al. 1990; Kleinman 1998). These pulsations change over monthly to yearly timescales, but the dominant modes tend to be around 600 s. Spectroscopic determination of its fundamental atmospheric parameters, i.e., $T_{\text{eff}}$ and $\log g$ (cgs), are uncertain due to the pulsations. For example, Koester et al. (2005) find $T_{\text{eff}} = 12,100$ K and $\log g = 7.9$, while (Bergeron et al. 2004) find $T_{\text{eff}} = 11,820$ K and $\log g = 8.14$. Accurate determination of the temperature and gravity is important in the case of G29-38 since diffusion calculations show orders of magnitude different settling times between 9000 and 14,000 K (see Fig. 2 in Koester & Wilken 2006). With the two different measurements, G29-38’s predicted settling timescale would vary between 4 and 7 days. However, further uncertainties may be present for pulsators. Empirical observations of equivalent width variations can directly probe diffusion calculations.

Von Hippel & Thompson (2007, hereafter vHT07) investigated the equivalent width variability of the strong Ca ii K line by compiling all spectroscopic observations of G29-38 from 1996 to 2000. This data included low-resolution data sets used to monitor variability in the line-of-sight velocities and line profiles of hydrogen lines arising from G29-38’s pulsations (van Kerkwijk et al. 2000; Thompson 2006). This variability is correlated with pulsations of the white dwarf and are shorter than the settling time for metals (Koester & Wilken 2006). VHT07 found a significant dip in the equivalent width of the Ca ii K line of G29-38 between 1997 and 1999, from averages of the two low-resolution data sets. The measured equivalent width went from 167 $\pm$ 4 to 280 $\pm$ 8 mA. They attributed this dip to a change in accretion rate, but vHT07 were cautious about the effect of the pulsations on variability of the Ca line strength. They used an analysis of their time series spectra to estimate the effect incomplete coverage pulsations may have on the line strength. They estimated that for observations of less than an hour, errors due to pulsations would be less than 5%. However, shorter exposure times were more sensitive to incomplete pulsation coverage; errors ranging from 8% for a 10 minute exposure to 13% for a 5 minute exposure were possible. Based on their time series analysis, they concluded that the Ca ii K line strength changes they observed were not due to pulsations. VHT07 subsequently combined the two low-resolution data sets with several other high-resolution spectroscopic observations spanning 1996 to 2000 and predicted that the Ca ii K line must be variable on timescales as short as 15 days.

To test for further variability, we observed G29-38 in two multiday runs separated by 4 months for a total of 19 individual observations with exposure times ranging from 100 to 500 s. All individual observations are short enough to determine the extent of possible systematic errors introduced by pulsations. When averaged together, each run covers enough time (~40 minutes) to average over potential pulsations and get an idea of the true average strength of the line. We therefore expected that our data in aggregate would be accurate to less than a few percent, but we expected that any single observation might vary by as much as 7%–13%. The observations, data analysis, and our conclusions are detailed in the following sections of this Letter.

2. OBSERVATIONS

We observed G29-38 with the blue chip of the MIKE spectrograph (Bernstein et al. 2003) installed at the 6.5 m Magellan
Clay Telescope at Las Campanas Observatory (Chile). Six spectra were collected on UT 2007 June 25–26 and another 13 spectra on UT 2007 October 24–27. Both runs used a 0.7” × 5.0” slit, yielding an average spectral resolution of \( R = 40000 \) at 3933 Å. The spectra cover wavelengths between 3335 and 5120 Å. Typically three observations were taken per night with the exception of October 24, where one 100 s exposure was taken and then three 200 s exposures. The exposure times in June were determined by poor (1”–1.3”) seeing and were between 300 and 500 s, whereas spectacular (≈0.5”) seeing conditions were present for the first two nights of the October run, followed by nominal seeing (≈0.8”) the last two nights, allowing shorter exposure times. The signal-to-noise ratio of each exposure is roughly 25, for each night is roughly 50, and for the averages between runs is about 80.

The data were extracted and flat-fielded using the MIKE reduction pipeline written by D. Kelson, with methodology described in Kelson et al. (2000) and Kelson (2003, 2007). All six individual spectra from the June run, 2400 s in total, and all 13 individual spectra from the October run, 2500 s in total, were combined into averaged spectra for each run to search for weak lines other than Ca. An example of the fidelity of these combined spectra is presented in Figure 1.

Absorption lines from Mg ii, Ca i, and Ca ii were detected. In particular, the Ca ii K line and the Mg ii λ4481 line fell on two different orders, a feature that we took advantage of to minimize systematic errors in the determination of the equivalent widths. Other lines from Mg i were detected at 3830, 3833, and 3838 Å. Ca i was detected at 4227 Å, and the Ca ii H line and another Ca ii line at 3737 Å were similarly detected.

In order to claim variability in a measured equivalent width, a careful examination of the measurement procedure and an accounting of systematic errors must be done. This step is important to robustly compare observations done at different telescopes (see § 3). To this end we developed a method for measuring the equivalent width of a line while ensuring that all the line optical depth was measured and that all systematic errors were accounted for. Inspection of the average spectra from June and October show that the line profile differs between the two nights and that the lines deviate from pure Gaussian shapes (see Fig. 1). The variable profile and slightly asymmetric wings may be caused by shifting velocity fields on the WD surface (Koester & Kompa 2007). We decided to directly integrate the line profile from the data, rather than fitting the line with a Gaussian and missing contributions to the equivalent width.

The continuum of some lines is affected by curvature introduced by the proximity of strong hydrogen Balmer lines. That curvature is removed by fitting a high-order polynomial. The choice of how to fit the continuum and the size of the continuum window to use can introduce systematic errors in the equivalent width. We used a polynomial fit to the continuum with a window of at least \( ± 5 \) Å on each side of the line center. To account for the systematic errors, we measured the equivalent width over all combinations of window size from \( ± 5 \) to \( ± 10 \) Å and polynomial fits ranging from third to twelfth order. We took the median measure of the equivalent width and the sum in quadrature of the measured standard deviation of measurements (the systematic error) and the estimated error due to signal-to-noise ratios.

An appropriate window around a line must be chosen to accurately measure the total equivalent width. We determined an optimal window by measuring the equivalent width for the averaged spectra from each run using a range of window sizes.

Once the continuum is fitted, the equivalent width will peak at a radial distance from the line center where the line no longer contributes. In the particular case of the Ca ii K line, we determined this distance to be \( ± 1.2–1.5 \) Å.

In addition to the Ca ii K line, we measured the equivalent width for some of the weaker lines we detected, Mg ii λ4481, Ca i λ4227, and the Ca ii H line. For these lines, we restricted ourselves to the averages of the two runs.

3. TESTING METAL LINE VARIABILITY

Figure 2 shows the equivalent width of the Ca ii K line from each exposure, in order of ascending sequential observation; i.e., the first spectrum taken in June is observation 1 and the last spectrum in October is observation 19. Each night has a specific symbol on the plot.

We find that the equivalent width of the Ca ii K line is stable over the time period we observed G29-38. The standard deviation of all 19 measurements is 20 mA, while the median estimated error of the measurements is 20 mA. These results suggest that if there is variability due to pulsations for these spectra it is <7%. This is less than a factor of 2 lower than that predicted by vHT07. We can then add observations over a small number of pulsation periods to increase the signal-to-noise ratio without fear of being subject to systematic errors.

Averaging the two runs gives equivalent widths of \( 262 ± 9 \) mA and \( 269 ± 7 \) mA for the June and October runs, respectively. These values are thus consistent with no variation in line strength over 4 months at the 2% level. If short-term variability occurred between the two measurements we would not have detected it. This implies that the underlying source of accretion is steady. Table 1 shows the results for our June and October runs (MJD 54278.5 and 54399.1, respectively) for all species measured.

It is relevant to revisit whether there has been significant variation of the equivalent width of the Ca ii K line or any of the other lines since 1997, the dates of the last published values. There exist eight other measurements, mentioned in vHT07,
Fig. 2.—Equivalent width of G29-38’s Ca ii K line from each observation in our 2007 June and October runs. The first six points are from the June run and exposure times range between 300 and 500 s in duration, while the last 13 points are from the October run. Exposure times in this case range between 100 and 200 s. The median of all the observations, 264 ± 20 mÅ, is represented by the horizontal solid line.

Fig. 3.—Equivalent width of G29-38’s Ca ii K line spanning the past decade. Squares are from high-resolution data, while diamonds are from low-resolution data. The horizontal line indicates the median of all the observations, not including the low-resolution data reported in vHT07.

TABLE 1

| MJD     | Ca ii K (mÅ) | Ca ii H (mÅ) | Ca i (mÅ) | Mg i (mÅ) |
|---------|--------------|--------------|-----------|-----------|
| 50636.577 | 268 ± 15     | 60 ± 16      | 24 ± 5    | 35 ± 11   |
| 51158.200 | 256 ± 9      | 48 ± 10      | 18 ± 5    | 47 ± 17   |
| 51403.463 | 248 ± 11     | 55 ± 5       | 23 ± 2    | 44 ± 7    |
| 51762.681 | 232 ± 12     | 57 ± 13      | ...       | ...       |
| 51804.670 | 239 ± 7      | 51 ± 6       | ...       | ...       |
| 54278.500 | 262 ± 9      | 50 ± 15      | 26 ± 4    | 51 ± 20   |
| 54399.100 | 269 ± 7      | 61 ± 9       | 23 ± 2    | 43 ± 7    |

Fig. 4.—Top: Plot of measured Mg ii λ4481 line equivalent width; horizontal line is median value of 44 ± 6 mÅ. Middle: Plot of measured Ca i λ4427 line equivalent width; horizontal line is median value of 23 ± 3 mÅ. Bottom: Plot of measured Ca ii H to K equivalent width ratios for G29-38; horizontal line is median value of 0.22 ± 0.02

We also measured the equivalent widths of the other three metal lines in the Keck and VLT data. The results, combined with our data, are shown in Figure 4. Similar to the strong Ca ii K line, these weaker lines detected in G29-38 show no variability. The Mg ii λ4481 line is visible in our data as well as the Keck data, but it is not reliably detected in the VLT data due to insufficient resolution. Over the five available observations this line is stable at a 12% level, with a median equivalent width including results from Koester et al. (1997), Zuckerman et al. (2003), and Koester & Wilken (2006). Excluding the low-resolution data from vHT07, we are left with six high-resolution spectra of G29-38 in addition to our own data. Of the six, three are from the Keck Observatory and are publicly available (see Zuckerman et al. 2003), while two spectra from the ESO VLT were kindly provided by D. Koester. We disregard the discovery equivalent width due to its large error.

We remeasured the equivalent width of the Ca ii K line using the same method described in § 2, and compared the new values to our data. Figure 3 shows the results, including the low-resolution data reported in vHT07. Neglecting the two low-resolution points, the median of the seven high-resolution observations is 262 mÅ, with a standard deviation of 7.6 mÅ and a median measurement error of 10.6 mÅ. With the exception of the low-resolution points, the underlying source of accretion for G29-38 is incredibly stable over a timescale of 10 years. If we add the reported measurements of the low-resolution data, the 1997 point remains significantly discrepant. We compare the 1999.65 point with the median and find that it is consistent at <2 σ. However, as noted in § 4, our measured equivalent widths are higher on average than vHT07’s and so the 1999.65 point may also be discrepant if we were to remeasure that data.
of 44 ± 6 mÅ. The Ca i λ4427 line was also detected in our data and the Keck data, but was not detected in the VLT data for the same reason as the Mg i line above. Over the five available observations, the Ca i line equivalent width is 23 ± 3 mÅ, and is stable at a 13% level. Finally, we measured the Ca ii H line, which resides in the wings of the He line. This line also has a stable equivalent width of 55 ± 5 mÅ over 10 years at a 10% level. This line was detected in all the available data. The results can be found in Table 1.

Given optically thin conditions for the Ca ii H and K lines, their strength should be in the ratio of H to K of 1 : 2. We plot the observed line ratios over the Keck, VLT, and Magellan data in the bottom diagram in Figure 4. The ratio is 22% ± 2%, which can be explained by additional opacity at the H line due to the He line wings (D. Koester 2007, personal communication).

4. CONCLUSIONS

We have shown that G29-38’s Ca ii K line has not significantly varied recently, or over the past decade, with the exception of one or possibly two observations. A stable result is also obtained for the other metal lines Mg ii λ4481, Ca i λ4427, and Ca ii H. We see no evidence for systematic errors due to pulsations at around the 7% level in our observations, slightly at odds with the predictions of vHT07. However, G29-38’s modes are variable in both period and strength, while vHT07 calculated errors for a worst case scenario involving strong pulsations. The last measurements of G29-38’s pulsations that we are aware of occurred in 2005 and the strongest modes had flux variations of ∼1%, whereas the flux variations reported for vHT07’s low-resolution data were ∼3% (M. Montgomery 2008, private communication). It is therefore not surprising that we do not see evidence for errors due to pulsations. Other calculations of equivalent width variability due to pulsations also predict errors at the level of 10%, suggesting that when accounted for, these errors have minimal impact on long term observations (Koester & Kompa 2007).

One point to note is the different values that we derive in our measurement of the archival equivalent widths when compared to those obtained by vHT07. Those differences come from the different techniques used in the measurement of the equivalent width and in the continuum fitting, which highlights the need for a consistent method when comparing observations from different sources. The fact that our method produces on average larger equivalent widths than those reported by vHT07 can be related to incorrect accounting for line flux in the wide wings of the Ca ii K line profile when using Gaussian fits.

Since the low-resolution data are the only points that deviate more than a few percent from the median equivalent width value, it is important to determine whether low-resolution observations could give spurious results at levels greater than their measured errors. The lack of any flat-fielding in vHT07, as well as changes in the wings of the He line, may all affect continuum fitting, something that low-resolution data will be more sensitive to. To test this hypothesis, we rebinned our data down to different resolutions that approach the 7–8 Å level used in vHT07 to see whether we could recover the same equivalent width using our method outlined above. At this level of rebinning, our S/N exceeded that reported in vHT07. We used a window for measuring the equivalent width of 10 Å centered on the line and fit the continuum with 10 Å on either side of the window. We tried resolutions that ranged from 0.04 to 4 Å. We find for the average spectrum of our October data that the measured equivalent widths at the lowest resolutions deviated significantly from our measured value at high resolution, going as low as 140 mÅ to as high as 350 mÅ. At resolutions less than 0.1 Å, our results converge to within 1 σ of our previously measured value. It is impossible to tell whether vHT07’s data suffers from a similar problem, but high-resolution, high-S/N confirmation of this variability is desired.

Whether the accretion onto G29-38 varies or not, the idea of variable accretion onto white dwarfs is an interesting line of study. Spectroscopic monitoring of DAZs with close companions and short settling times offers an interesting insight into the stellar wind behavior of M dwarfs (Debes 2006), while variability in single DAZs may conceivably determine the source of material that pollutes the surface of these white dwarfs. Both types of polluted white dwarfs directly test models of white dwarf diffusion times. Few DAZs have been monitored in a systematic way, but the MIKE spectrograph provides a sensitive and stable platform with which to do long-term studies.

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