Compressing large data sets to a manageable number of summaries that are informative about the underlying parameters vastly simplifies both frequentist and Bayesian inference. When only simulations are available, these summaries are typically chosen heuristically, so they may inadvertently miss important information. We introduce a simulation-based machine learning technique that trains artificial neural networks to find non-linear functionals of data that maximise Fisher information: information maximising neural networks (IMNNs). In test cases where the posterior can be derived exactly, likelihood-free inference based on automatically derived IMNN summaries produces nearly exact posteriors, showing that these summaries are good approximations to sufficient statistics. In a series of numerical examples of increasing complexity and astrophysical relevance we show that IMNNs are robustly capable of automatically finding optimal, non-linear summaries of the data even in cases where linear compression fails: inferring the variance of Gaussian signal in the presence of noise; inferring cosmological parameters from mock simulations of the Lyman-$\alpha$ forest in quasar spectra; and inferring frequency-domain parameters from LISA-like detections of gravitational waveforms. In this final case, the IMNN summary outperforms linear data compression by avoiding the introduction of spurious likelihood maxima. We anticipate that the automatic physical inference method described in this paper will be essential to obtain both accurate and precise cosmological parameter estimates from complex and large astronomical data sets, including those from LSST and Euclid.

Current data analysis techniques in astronomy and cosmology often involve reducing large data sets into a collection of sufficient statistics (Bond et al. 1998, Heavens et al. 2000, Tegmark et al. 1997). There are several methods for condensing raw data to a set of summaries. Amongst others, these methods could be: principal component analysis (PCA) (Connolly et al. 1995, Francis et al. 1992, Lahav 2009, Madgwick et al. 2002, Murtagh and Heck 1987); statistics including the mean, covariance, and higher point functions (Belmon, L. et al. 2002, Betancort-Rijo 2012) or; calculating the autocorrelation or power spectrum (Betancort-Rijo 2012, Segal 2012). Unfortunately, summaries calculated using the above methods can still be infeasibly large for data-space comparison. For example, analysis of weak lensing data from the Euclid and the Large Synoptic Survey Telescope (LSST) photometric surveys will have around $10^4$ summary statistics (Heavens et al. 2017). Reducing the number of summaries further results in enormous losses in the information available in the raw data (Heavens et al. 2017).

Another popular way of summarising data is using the Massively Optimised Parameter Estimation and Data (MOPED) compression algorithm (Heavens et al. 2000). Summaries from MOPED are linear combinations of data that compress the number of data points down to the number of parameters of a model describing the data. MOPED is completely lossless when noise in the data is independent of the parameters and when the likelihood is, at least to first order, Gaussian (Heavens et al. 2000). The MOPED algorithm has been used on many problems in astronomy and cosmology such as studying the star formation histories of galaxies (Heavens et al. 2004, Panter et al. 2007, Reichardt et al. 2001), analysing the cosmic microwave background (Gupta and Heavens 2002, Zablocki and Dodelson 2016), and identifying transients (Protopapas et al. 2005) to name but a few. Unfortunately, using linear combinations of the data for compression may not be optimal for maximising the possible information available, even when the likelihood is known (Alsing and Wandelt 2017).
For many astronomical and cosmological problems, it can become impossibly difficult to write a likelihood function which describes, not only physics, but also includes any selection bias and instrumental effects. Recently, methods have become available to perform inference when a likelihood is not available via approximate Bayesian computation (ABC). ABC is a technique which allows samples to be drawn from an approximate posterior distribution. Forward simulations are first created using parameter values drawn from a prior and samples are accepted or rejected by comparing the distance of the simulation to the real data. To efficiently approach the true posterior distribution, it is convenient to couple ABC with a sampling procedure such as population Monte Carlo (PMC). ABC using PMC (PMC-ABC) is a method to obtain approximate parameter distributions by iterating through weighted samples from the prior [Pritchard et al. 1999, Tavare et al. 1997] and can massively reduce the number of samples which need to be drawn during ABC.

Likelihood-free inference has been used for a variety of astronomical problems which include deducing quasar luminosity functions (Schafer and Freeman 2012), understanding early time galaxy merger rate evolution (Cameron and Pettitt 2012), constraining cosmological parameters with supernova observations (Weyant et al. 2013), interpreting galaxy formation (Robin et al. 2014), searching for the connection between galaxies and halos (Hahn et al. 2017), measuring cosmological redshift distributions (Kacprzak et al. 2017), inferring photometric evolution of galaxies (Carassou et al. 2017), and calculating the ionising background using the Lyman-α and Lyman-β forest transmission (Davies et al. 2017). Each of the above examples are used in conjunction with publicly available (PMC-)ABC codes (Akeret et al. 2015, Ishida et al. 2015, Jennings et al. 2016).

A two-step compression algorithm was defined in (Alsing et al. 2018) that is capable of optimally summarising data whilst preserving information when the likelihood is not known. The first step involves extracting informative statistics from raw data (or simulations of the data) heuristically, i.e. perhaps using the power spectrum or using PCA. The summaries of the simulations contain information about physics, selection bias and the instrument. A second step then assumes an asymptotic likelihood to perform compression from the summaries gathered in the first step down to the number of parameters in the model as in MOPED or (Alsing and Wandelt 2017). The choice of likelihood in the second step does not bias the inference of model parameters during ABC, although the compression will be closer to optimal by choosing a better likelihood function.

However, what if there is information in the data that we did not think to summarise in a first-step summary? In this paper we introduce the concept of information maximising neural networks (IMNNs). Through the use of machine learning, we can circumvent the two step compression used in (Alsing et al. 2018) and find the most informative non-linear data summaries by training a neural network using the Fisher information matrix as a reward function. In fact, if we already know some informative summaries, such as those calculated in the first step of (Alsing et al. 2018), we can use the IMNN to calculate summaries of the data which optimally increase the information further and then including the IMNN summaries amongst the first-step summaries.

Once the network is trained, ABC proceeds as before. Model parameters can be drawn from a prior, used to generate simulations and once they are fed through the network, the IMNN summaries of the simulation can be compared to the summaries of the real data. Samples can then be accepted or rejected given the distance of the network summary of the simulation to the network summary of the real data to build the approximate posterior distribution of model parameters. The IMNN provides a framework to perform automatic physical inference simply by producing simulations.

In section I we describe how to calculate the Fisher information matrix and how linear summaries of the data can conserve Fisher information using the MOPED algorithm. In section II we lay out the procedure for creating non-linear summaries of the data. An overview of how artificial neural networks work is presented in section III and we continue in section IV by showing how maximising the determinant of the Fisher information matrix allows a network to be trained to provide the optimal non-linear set of summaries. Next, in section V we trace the steps to obtain parameter constraints from PMC-ABC using the network trained as prescribed in section IV. Finally, in section VI we give some test examples. The first test model provides an example where a single linear summary of the data would provide nearly no information about a parameter, but the non-linear summary provided by a trained artificial neural network can extract the maximum information the data contains. The second example is more astronomically motivated, using the absorption of flux from quasars by neutral hydrogen to constrain the amplitude of scalar perturbations. Finally we use the network to summarise and constrain the central oscillation frequency of a gravitational wave burst from Laser Interferometer Space Antenna (LISA). This problem was used in (Graff et al. 2011) to show that MOPED compression introduces spurious maxima in the posterior distribution; we show that the non-linear IMNN data compression introduced in this paper can avoid this peculiarity.

I. FISHER INFORMATION AND LINEAR COMPRESSION

A likelihood function \( L(d|\theta) \) of some data, \( d \), with \( n_d \) data points, is informative about a model with a set of \( n_\theta \) parameters, \( \theta \). The more sharply peaked \( L(d|\theta) \) is at a particular value of \( \theta \), the better \( \theta \) is known. The Fisher
information describes how much information \( d \) contains about the linear parameters, \( \varphi \), and can be calculated by finding the second moment of the score of the likelihood (Fisher 1925, Kendall and Stuart 1969, Kenney and Keeping 1951), i.e. the variance of the partial derivative of the natural logarithm of the likelihood with respect to the parameters at a fiducial parameter value, \( \varphi^{(d)} \).

The MOPED compression is lossless in the near combination of data which optimises the linearised algorithm (Heavens et al. 2000), parameters of the model, the number of points in the data, convenience.

Given parameters \( d \), where \( d \) is the data and \( \mu(\varphi) \) is the mean of the model given parameters \( \varphi \), which we will denote \( \mu \) for convenience. \( C \) is the covariance of the data and is assumed to be independent of the parameters. Using the MOPED algorithm (Heavens et al. 2000), \( d \) can be compressed from the number of points in the data, \( n_{data} \), to the number of parameters of the model, \( n_{\varphi} \), simply by seeking the linear combination of data which optimises the linearised parameters. The MOPED compression is lossless in the sense that the Fisher information is conserved under the transformation

\[
x_{\alpha} = r_{\alpha}^T d
\]

where \( \alpha \) labels the parameter and \( r_{\alpha} \) is calculated by maximising the Fisher information ensuring that \( r_{\alpha} \) is orthogonal to \( r_{\beta} \) (where \( \alpha \neq \beta \)). The form of \( r_{\alpha} \) is

\[
r_{\alpha} = \frac{C^{-1} \mu_{\alpha}}{\sqrt{\mu_{\alpha}^T C^{-1} \mu_{\alpha}}},
\]

for the first parameter, \( \varphi_1 \), and where \( \partial/\partial \varphi_\alpha \equiv \mu_\alpha \). For each parameter afterwards,

\[
r_{\alpha} = \frac{C^{-1} \mu_{\alpha} - \sum_{i=1}^{n-1} (\mu_{\alpha}^T r_i) r_i}{\sqrt{\mu_{\alpha}^T C^{-1} \mu_{\alpha} - \sum_{i=1}^{n-1} (\mu_{\alpha}^T r_i)^2}}.
\]

After creating the linear summaries, \( x = \{x_{\alpha} | \alpha \in [1, n_{\varphi}] \} \), \( x \) is as informative about \( \varphi \) as \( d \) is with regards to the Fisher information, for the likelihood in equation (1.4). The Fisher information takes the form

\[
F_{\alpha \beta} = \text{Tr} [\mu_{\alpha}^T C^{-1} \mu_{\beta}]
\]

The lossless compression of the data, \( d \rightarrow x \), is only possible when the likelihood is exactly of the form in equation (1.4). Nearly lossless compression is still possible if the peak of the likelihood is approximately Gaussian. Often, this will be a good approximation in the asymptotic limit, i.e., when the data are informative about the parameters.

II. NON-LINEAR FISHER INFORMATION
MAXIMISING SUMMARIES

We are influenced by the MOPED algorithm to find some transformation which maps the data to compressed summaries, \( f : d \rightarrow x \), whilst conserving Fisher information, but without the limitation that the method is only valid as a Gaussian approximation. \( f \) is a function that modifies the original likelihood describing the data, which need not be known \( a \) priori, into the form

\[
-2 \ln \mathcal{L}(x|\varphi) = (x - \mu_f(\varphi))^T C_f^{-1} (x - \mu_f(\varphi))
\]

where

\[
\mu_f(\varphi) = \frac{1}{n_\alpha} \sum_{i=1}^{n_\alpha} x_i^f,
\]

is the mean value of \( n_\alpha \) summaries, \( \{x_i^f | i \in [1, n_\alpha] \} \), where each summary is obtained from a simulation \( d_i^f = f(\varphi, i) \) using \( f : d_i^f \rightarrow x_i^f \). We will denote \( \mu_f(\varphi) \equiv \mu_f \) for convenience. Each \( i \) denotes a different random initialisation of a simulation. Similarly \( C_f^{-1} \) is the inverse of the covariance matrix which is again obtained from simulations of the data

\[
(C_f)_{\alpha \beta} = \frac{1}{n_{\alpha} - 1} \sum_{i=1}^{n_{\alpha}} (x_i^f - \mu_f)_\alpha (x_i^f - \mu_f)_\beta.
\]
Using equation (1.2) a modified Fisher information matrix can be calculated from the likelihood in equation (2.1)

\[
F_{\alpha\beta} = \text{Tr} \left[ \mu_{f;\alpha}^T C_f^{-1} \mu_{f;\beta} \right].
\] (2.4)

Here, the values of \( \mu_{f;\alpha} \) and \( C_f^{-1} \) are calculated using fixed, fiducial parameter values, \( \Theta_{\text{fid}} \), such that the simulations are \( d_{f;\text{fid}} = d(f(\Theta_{\text{fid}}), i) \). Although \( f : d \to x \) is not specified, a subclass of \( f \) is accessible via a neural network, described in detail in section [III]. We will show how this function can be found by training a neural network in section [IV].

### III. ARTIFICIAL NEURAL NETWORKS

Artificial neural networks are arbitrary maps from some inputs to outputs. Consider some data vector \( d = \{ d_i | i \in [1, n_d] \} \) with \( n_d \) data points. Each data point is regarded as an input to a network. For a deep neural network, a series of hidden layers are able to learn levels of abstraction from the input [Bengio 2009, Cybenko 1989, Deng and Yu 2014, Goodfellow et al. 2016, Nielsen 2015]. Each layer, \( l \), of the network contains a set of neurons which takes some number of inputs and provides one output per neuron [McCulloch and Pitts 1943, Pitts and McCulloch 1947]. The output a neuron is activated by a non-linear activation function

\[
a^l_i = \phi \left( v^l_i \right)
\] (3.1)

where

\[
v^l_j = \sum_i w^l_ji a^{l-1}_i + b^l_j,
\] (3.2)

is a weighted, biased input at each layer with weights \( w^l_j \equiv w^l_{ji} \) and biases \( b^l_j \equiv b^l_{ji} \) [McCulloch and Pitts 1943]. \( i \) describes an element of the output vector of a collection of neurons in the \( (l-1) \)th layer and \( j \) indexes the neuron in layer \( l \). With these notations, the input to the network can be considered to be the output of a zeroth layer of a network, \( d_l \equiv a^0_L \). Stacking several neurons into a hidden layer and stacking several hidden layers, taking the outputs from the previous layer as the inputs to each node in the next layer, allows for greater levels of abstraction from the input data [Deng and Yu 2014]. These networks are often referred to as deep networks. Note that the addition of too many layers can lead to expensive computations and overfitting by the network so that it becomes difficult to train. The network output at the final layer can be described by \( a^L = \{ a^L_i | i \in [1, n_{\text{outputs}}] \} \) where \( n_{\text{outputs}} \) is the number of outputs in the final layer, labelled \( L \), and \( a^L_i = \phi (v^L_i) \).

As mentioned at the end of section [III], a neural network can be used as a representation of \( f : d \to x \), which compresses data to summary statistics. Formally, this subclass of functions is described, for some input \( z \), by

\[
f^l : z \to a^l = \phi \left( \sum_i w^l_{ji} f^{l-1}(z)_i + b^l_j \right),
\] (3.3)

for \( l > 0 \) and

\[
f^0 : z \to a^0 = m,
\] (3.4)

where the compressed summary is given at \( l = L \) of the recursion and the input to the function at \( l = 0 \) is taken to be the identity.

#### 1. Activation functions

The activation function, \( \phi (v^l_j) \), in equation (3.1) describes whether the artificial neuron fires or not, i.e., whether the inputs are informative or useful for describing the output [Cybenko 1989, He et al. 2015, Krizhevsky et al. 2012, Nielsen 2015]. It is the activation function that provides the non-linearity necessary for the neural network to learn the complex map from inputs to outputs by combining the relevant combinations of inputs at each layer in a non-trivial way. As long as there are enough hidden layers, the form of the activation function is relatively unimportant since the weights and biases will be trained to combine the outputs of each hidden layer in such a way as to provide the correct map. There are many options for the choice of activation function, including tanh and sigmoid functions. Currently popular activation functions are the rectified linear unit (ReLU) [He et al. 2015]. We show here, as an example, an adaptation called leaky ReLU

\[
\phi(x) = \begin{cases} 
\alpha x & x \leq 0 \\
0 & x > 0 
\end{cases}
\] (3.5)

where \( \alpha = 0 \) for ReLU and \( \alpha \) is small and positive for leaky ReLU [Maas et al. 2013]. Although the ReLU family of activation functions are linear, stacking several layers of neurons provides a function which approximates a non-linear function, and is extremely quick to calculate. It will become apparent that the derivative of the activated output with respect to the weighted, biased inputs are essential for training neural networks. The derivative of the ReLU family of activation functions can also be efficiently calculated as

\[
\frac{\partial \phi(x)}{\partial x} = \begin{cases} 
\alpha & x \leq 0 \\
1 & x > 0 
\end{cases}
\] (3.6)

Although we have shown ReLU as an example, we explore various activation functions across the population of networks that we train.

#### 2. Back propagation

A scalar loss function, \( \Lambda (a^L) \), is calculated from the outputs of the network \( a^L \). In supervised deep learning,
the loss function describes how far the outputs are from a set of labels for the training data (Rumelhart et al. 1986). An iterative procedure, called back propagation, uses the chain rule to find how much the weights and biases need to change to minimise the loss function (Rumelhart et al. 1986). Using gradient descent (Kiwiel 2001) it can be seen that the weights and biases must be updated using

\[ w_{ji} \rightarrow w_{ji} - \eta \frac{\partial \Lambda}{\partial w_{ji}} \]  

(3.7)

and

\[ b_i \rightarrow b_i - \eta \frac{\partial \Lambda}{\partial b_i} \]  

(3.8)

where \( \eta \) is a tunable learning rate which dictates the size of the steps that the weights and biases are able to take on each update (Nielsen 2015). It is very efficient to calculate the derivatives in equations (3.7) and (3.8) at the last layer using

\[ \frac{\partial \Lambda}{\partial v_i^l} = \frac{\partial \Lambda}{\partial a_i^l} \frac{\partial a_i^l}{\partial v_i^l} \]  

(3.9)

From any layer, the rate of change of the loss function with respect to the weighted, biased inputs at the previous layer can be found using

\[ \frac{\partial \Lambda}{\partial v_i^l} = \sum_j u_j^{l+1} \frac{\partial \Lambda}{\partial a_i^l} \frac{\partial a_i^l}{\partial v_i^l} \]  

(3.10)

The changes in the loss function under changes in the weights or the biases are then calculated using

\[ \frac{\partial \Lambda}{\partial w_{ji}} = \frac{\partial \Lambda}{\partial v_j^l} \frac{\partial v_j^l}{\partial w_{ji}} = \frac{\partial \Lambda}{\partial v_j^l} a_i^{l-1} \]  

(3.11)

and

\[ \frac{\partial \Lambda}{\partial b_i^l} = \frac{\partial \Lambda}{\partial v_i^l} \frac{\partial v_i^l}{\partial b_i^l} = \frac{\partial \Lambda}{\partial v_i^l} \]  

(3.12)

Each of the \( a_i^l \) and the derivatives with respect to the weighted biased inputs (using equation (3.6)) are calculated on the forward pass of the network inputs. Back propagation allows the change of the loss function with respect to all of the weights or biases to be calculated in just one pass forward and one pass backwards (Nielsen 2015). By calculating the change in the loss function with respect to the network outputs and applying equation (3.9) successively, the weight and bias updates at every layer can be calculated easily.

The back propagation procedure is repeated many times using different sets of training inputs (Nielsen 2015). Once all of the training inputs are used, one epoch of training is complete. After one epoch of training, the order of the training inputs can be jumbled and the training procedure repeated many times until the loss function is minimised (Nielsen 2015).

3. Overfitting

It is possible that the network weights become tuned to features in the training data which are not present in the real data. To prevent this overfitting, we implement dropout (Srivastava et al. 2014). Dropout is a technique where a random fraction of the neurons are set to zero on each batch of training and after back propagation only the weights and biases of the active neurons are updated. Performing dropout during training equates to training many sub-networks, where all the neurons share weights and biases. Each of the sub-networks can learn specific features in the data, but the consensus network does not learn features too strongly.

4. Training and test data sets

When training a network, it is essential to test how well the network is learning by using a test set which contains data which is not present in the training set. However, it is extremely important to note that even the accuracy of prediction on the test set should not be considered to be a measure of the predictive ability of the network. It is considered normal to tune a network to achieve the minimal loss of the test set without showing signs of overfitting. A third, completely unseen, data set should then be used to quote network accuracies. In doing so, the irreproducible accuracy scores often quoted in the literature, arising from only considering a network that is highly tuned on the test set, are avoided. In this paper we train and test networks with a training set and a test set and use the comparison between the posterior distribution obtained using the network output and the analytically calculated distribution as our confirmation that the network is accurate.

IV. FINDING NON-LINEAR SUMMARIES

Inspired by supervised artificial neural networks we are able to create a network capable of maximising the Fisher information to create non-linear summaries of data. The output of the network, \( \mathbf{x} \equiv \mathbf{a}^L \) is a compressed summary of some data, \( \mathbf{d} \). Since the data is a function of some parameters, \( \mathbf{\theta} \), given some model, the summary can be described as a function of these parameters, as well as the weights and biases at each layer, \( l \), of a network, \( \mathbf{x} \rightarrow \mathbf{x} (\mathbf{\theta}, \mathbf{a}^l, \mathbf{b}^l) \). The mean, \( \mu_l \), covariance, \( \Sigma_l \) and Fisher information matrix, \( \mathbf{F}_{\mathbf{\alpha} \mathbf{\beta}} \), from equations (2.2), (2.3) and (2.4), each become functions of
the weights and biases as well. Summaries of simulations, $x_i^a$, are obtained by passing simulations, $d_i^a$, through the network $f : d_i^a \rightarrow x_i^a$.

To compute the Fisher information matrix in equation (2.3), the derivative of the network needs to be calculated with respect to the parameters at fiducial values. It is, in principle, simple to find the derivative of the network with respect to the parameters due to partial derivatives commuting with sums

$$\mu_{i,\alpha} = \frac{\partial}{\partial \theta_{\alpha}} \frac{1}{n_s} \sum_{i=1}^{n_s} x_i^s |_{\theta_{\alpha}=\theta_{\alpha}^f}$$

$$= \frac{1}{n_s} \sum_{i=1}^{n_s} \left( \frac{\partial x_i}{\partial \theta_{\alpha}} \right) s_i^f |_{\theta_{\alpha}=\theta_{\alpha}^f}. \tag{4.1}$$

Unfortunately, since the parameters only appear in the simulations, numerical differentiation needs to be performed. The numerical differentiation is achieved by producing three copies of the simulation, $d_i^{x-s} = d_i^x (\hat{\theta}^x - \Delta \theta^+, i)$, and $d_i^{x-s} = d_i^x (\hat{\theta}^x + \Delta \theta^-, i)$, where $\Delta \theta^\pm$ is some small deviation from the fiducial parameter value. The derivative of the network output with respect to the parameters is therefore given by

$$\left( \frac{\partial x_i}{\partial \theta_{\alpha}} \right) s_i^f \approx x_i^{s, fid} - x_i^s |_{\theta_{\alpha}=\theta_{\alpha}^f} \frac{\partial \Delta \theta^+}{\partial \theta_{\alpha}} - x_i^s |_{\theta_{\alpha}=\theta_{\alpha}^f} \frac{\partial \Delta \theta^-}{\partial \theta_{\alpha}}. \tag{4.2}$$

Setting the random seed, $i$, to the same value when generating $d_i^{x-s}$ and $d_i^{x-s}$ suppresses the sample variance in estimates of the derivative of the mean. Although the network output can vary a lot between different simulations, the derivative with respect to parameters is much more stable to changes in the parameter value, meaning relatively few extra simulations ($n_{\theta|\theta} < n_s$) need to be computed to calculate the gradient of the mean.

Another way of calculating the derivative of the mean of the network output is to calculate the adjoint gradient of the simulations, and calculate the derivative of the network with respect to the simulations

$$\mu_{i,\alpha} = \frac{1}{n_s} \sum_{i=1}^{n_s} \sum_{k=1}^{n_k} \frac{\partial x_i^{s, fid}}{\partial \hat{d}_k} \frac{\partial \hat{d}_k^{s, fid}}{\partial \theta_{\alpha}}, \tag{4.3}$$

where $i$ labels the random initialisation of the simulation and $k$ labels the data point in the simulation. In certain situations, calculating the adjoint gradient of the simulations may be more efficient than the method described in equations (4.1) and (4.2).

One simple way of obtaining the optimal non-linear summary from some data is to maximise the determinant of the Fisher information matrix calculated from the network, $|F|$

$$\Lambda = -\frac{1}{2} |F|^2. \tag{4.4}$$

The Fisher information matrix terms are produced from the second derivatives of the Kullback-Leibler divergence, i.e. the information gain, and is hence directly related to the Shannon entropy $\text{Kullback1968}$. In particular, the Fisher information matrix is the Shannon information of the Gaussian probability distribution function which optimally approximates the likelihood in (2.1) near its peak. For this reason, choosing to maximise the determinant of the Fisher information is equivalent to maximising the Shannon information of this distribution. The error is then found by taking the derivative of the loss function with respect to the network output. Normally $x_i^a$ would be considered as the network output when the input is a simulation, but since the quantity of interest in our problem is statistically calculated over a large number of network outputs, we follow the cartoon in figure 1 and use the determinant of the Fisher information matrix as the true network output. First, a large number of simulations at fixed fiducial parameter value and random initialisation (as well as the simulations created to calculate the derivative of the mean) are fed forwards through identical networks. All the network outputs from the fixed fiducial parameter simulations are used to calculate the covariance as in equation (2.3). Meanwhile, the rest of the network outputs are used to find the derivative of the mean with respect to the parameter as in equations (4.1) and (4.2). These are combined to give the Fisher information matrix of equation (2.4). If we consider the true network output to be $a^L = |F|$ rather than $x^a$ then the error can be defined as

$$\frac{\partial \Lambda}{\partial a^L} = -|F|. \tag{4.5}$$

Training then commences over many epochs of weight and bias updates until the Fisher information stops increasing. In practice, a problem arises when using equation (4.5), since the Fisher information is invariant under linear scaling of the summary. To control the magnitude of the summaries we can artificially induce a scale by adding the determinant of the covariance matrix, $|C_i|$, to the error function

$$\frac{\partial \Lambda}{\partial a^L} = -|F| + |C_i|. \tag{4.6}$$

The network is penalised when the determinant of the covariance is large. When using equation (4.6) the network provides the summary which maximises the Fisher information whilst minimising the covariance of the outputs.

Although the network is capable of extracting all necessary summaries of the data without any prior knowledge of what the parameters represent, we can imagine the IMNNs would be better suited to extending the heuristic first-step summaries. For example, if the power spectrum is a known useful summary of some data, the network can be trained to find any statistic which increases the Fisher information further. With the power spectrum and the
network summary, a second stage compression as described in (Alsing and Wandelt 2017) can be used for efficient parameter inference. This way, inexhausted information of the data can be unlocked, even when the form of the data combination that probes it is not known.

V. APPROXIMATE BAYESIAN COMPUTATION

Approximate Bayesian computation (ABC) is a technique of finding an approximate posterior distribution for some model parameters by accepting or rejecting samples dependent on how similar simulations created using the sample parameters are to the real data (Rubin 1984]. It is useful to choose an appropriate sampling procedure to quickly approach the true posterior for the parameters without creating too many simulations. Population Monte Carlo (PMC) is an algorithm by which samples can be obtained by iterating through weighted draws from a prior, even when the likelihood is not accessible (Kitagawa 1996). Although PMC has a variety of uses, such as filtering, we are going to couple it to ABC (PMC-ABC) to effectively approach the true posterior (Pritchard et al. 1999, Tavaré et al. 1997).

ABC algorithm starts by drawing $\theta^0_k$ from prior, even when the likelihood is not accessible. Iterations can be performed until the number of draws is much larger than the number of wanted samples from the posterior, $N$. A large number of draws compared to the number of accepted parameter values is a sign that the approximate posterior has stopped changing considerably between iterations.

VI. TESTING INFERENCE WITH INFORMATION MAXIMISING NEURAL NETWORKS

In this section we use the information maximising neural network on a range of test models. In section VIA we use the network to summarise a Gaussian signal with unknown variance, as well as Gaussian signal with unknown variance that was contaminated by noise, first of known variance and then of unknown variance. We consider the same problem in section VIB showing that the network provides nearly optimal, informative summaries in spite of a poorly chosen fiducial parameter value by learning the correct map. In section VIC we constrain the amplitude of scalar perturbations using simulations of quasar absorption spectra which can be summarised by a single statistic provided by the network. Finally, in section VID we demonstrate the performance of IMNN compression for the case estimating the central frequency of a LISA gravitational wave chirp. This example addresses a concern raised in (Graff et al. 2011) where the authors show that a linear summary of data in the time
domain can be misleading about a parameter in the frequency domain. We are show that the non-linear summary avoids this problem and is more informative.

A. Summarising Gaussian signals

A simple toy model can be constructed where linear combinations of the data are unable to provide information about parameters.

Consider an experiment which measures \( n_d = 10 \) data points which are drawn from a zero-mean Gaussian where the variance, \( \vartheta = \sigma^2 \), is not perfectly known, \( \mathbf{d} = \{ d_i \mid i \in [1, n_d] \} \). The likelihood is written

\[
\mathcal{L}(\mathbf{d}|\vartheta) = \prod_{i=1}^{n_d} \frac{1}{\sqrt{2\pi\vartheta}} \exp \left[ -\frac{1}{2\vartheta} d_i^2 \right] = \frac{1}{(2\pi\vartheta)^{n_d/2}} \exp \left[ -\frac{1}{2\vartheta} \sum_{i=1}^{n_d} d_i^2 \right],
\]

such that

\[
-2 \ln \mathcal{L}(\mathbf{d}|\vartheta) = \frac{1}{\vartheta} \sum_{i=1}^{n_d} d_i^2 + n_d \ln [2\pi\vartheta].
\]
Fisher information is a minimal sufficient statistic. Maximising the (logarithm of the) likelihood with respect to the variance relates the value of the statistic to the variance
\[
\frac{\partial \ln L(x|\theta)}{\partial \theta} = \frac{x}{2\theta^2} - \frac{n_d}{2\theta} = 0
\]
so that
\[
x = n_d \bar{\theta}.
\]
The Fisher information is calculated using equation (6.2)
\[
F = \frac{x}{(\bar{\theta}^* d)^2} - \frac{n_d}{2(\bar{\theta}^* d)^2} = \frac{n_d}{2(\bar{\theta}^* d)^2}.
\]
For \(n_d = 10\) and a fiducial variance of \(\bar{\theta}^* d = 1\) the Fisher information is
\[
F = 5.\tag{6.7}
\]
Since the single summary is a non-linear combination (squared sum) of the data, linear combinations will not be able to provide a single sufficient statistic.

Now consider training a network to maximise the Fisher information whilst summarising the data, as laid out in section IV. We show the progress an example network makes until it extracts the full information in figure 2. The fully connected network has two hidden layers with 256 neurons in each. We denote this configuration makes until it extracts the full information in figure 2. We have found that a very large variety of hyperparameters will provide us with approximately \(F = 5\). Most notably we can use very deep networks with few neurons such as \([5, 5, 5, 5, 5]\) to extremely simple networks with large numbers of neurons, i.e. \([2048, 2048]\), each with very similar outcomes. The main difference with different architectures, that we have found, is the number of epochs necessary to maximise the Fisher information matrix. We have chosen a simple network of \([256, 256]\) since it seems to converge more quickly than other networks.

From equation (6.7), it can be seen that the maximum Fisher information attainable for this problem is \(F = 5\). Figure 2 shows that \(F = 5.15 \pm 0.39\) is obtained by the network over the last 10% of the training epochs. The solid blue line in figure 2 is the value of the Fisher information obtained by running a set of 500 simulations (and 50 partial derivatives), which are contained in the training set, through the network. The dashed orange line shows the Fisher information obtained by running the same number of simulations through the network, but where none of the simulations are present in the training set. The maximum amount of Fisher information expected is \(F = 5\), show as a black dashed line. It is clear that the network manages to extract the entirety of the information given the data.

We have found that a very large variety of hyperparameters will provide us with approximately \(F = 5\). Most notably we can use very deep networks with few neurons such as \([5, 5, 5, 5, 5]\) to extremely simple networks with large numbers of neurons, i.e. \([2048, 2048]\), each with very similar outcomes. The main difference with different architectures, that we have found, is the number of epochs necessary to maximise the Fisher information matrix. We have chosen a simple network of \([256, 256]\) since it seems to converge more quickly than other networks.

Since the test model can be written down analytically, the true posterior distribution for some simulated test data \(d\) (shown in table I) can be found and is plot-
Data | Value  
---|---  
\(d_1\) & \(-0.91903399\)  
\(d_2\) & \(-0.37322515\)  
\(d_3\) & \(-0.05613342\)  
\(d_4\) & \(1.20816746\)  
\(d_5\) & \(0.07649269\)  
\(d_6\) & \(-0.47171141\)  
\(d_7\) & \(-1.4756571\)  
\(d_8\) & \(-0.62946463\)  
\(d_9\) & \(-1.30334079\)  
\(d_{10}\) & \(-0.41441639\)

Table I: Values of the input parameters for the original simulated true data set where the data is Gaussian noise \(d = \{ d_i \cap N (0, 1) \mid i \in [1, n_d] \}\).

It is interesting to see the network outputs as a function of \(\theta\), without using the PMC procedure. By performing ABC by randomly drawing from the whole prior, and not honing in on the true distribution, we can plot the network output as a function of the variance drawn from

\[ \Delta \theta \text{ is the solid orange curve in figure 3.} \]

The prior distribution used here is uniform between \(\theta = (0, 10]\). A first approximation of the posterior distribution using the network, without creating any additional simulations can be found using the asymptotic likelihood by expanding equation (2.1) about the fiducial variance with \(\Delta \theta = (-1, 9]\). The asymptotic likelihood result is plotted in figure 3 in dashed blue. It can be seen that the peak of the posterior found using the asymptotic likelihood corresponds with the peak of the analytic posterior, although as expected the rest of the distribution quickly deviates from the analytic result. To perform PMC-ABC, \(N = 1000\) parameter values, \(\{ \theta_k^0 \mid k \in [1, N] \}\), are drawn from the uniform prior distribution, \(p(\theta)\), between \(\theta = (0, 10]\). The PMC procedure, described above, is then carried out to obtain 1000 samples from the approximate posterior. Using a criterion that there needs to be 2000 draws of \(\theta_k^0\) in iteration \(T\) to be convinced that the approximate posterior has converged requires a total of 10232 simulations. The width of the acceptance parameter is \(\epsilon^T = 0.086\) at the last iteration, \(T\), meaning that the network summary of each of the accepted network summaries are within a band of \(x_{ik}^T = x \pm 0.086\) of the network summary of the real data, \(x\). The histogram of the accepted points are shown in figure 3 in purple. The PMC-ABC posterior distribution follows the analytic posterior distribution exactly, showing that the network has successfully learned how to summarise the data.

1. This approximation is only true for \(\text{Abs}[\Delta \theta] \ll 1\). The interval chosen here is used only for plotting purposes.
the prior, shown in figure 4. The green points show the rejected samples and the purple points (under the black dashed line) show the accepted draws. The black dashed line shows the network output of the real data. There is a strong correlation between the network summary of the simulations and the value of $\vartheta$ used to create the simulation. Requiring that there are 1000 samples whose summaries are within $\vartheta_{ik}^T = x \pm \varepsilon^T$, where $\varepsilon^T = 0.086$, necessitates more than 600,000 draws from the prior, 50 times more draws than the PMC needs. It should be noted that the network summary is not equal to the value of $\vartheta$ and, in general can vary a lot by changing the network architecture, the initialisation of the weights or even just changing the order of the simulations used to train the network. The variation in the network summary is a manifestation of how the Fisher information is invariant under linear scalings of a sufficient statistic, although the scale of the statistic is able to be constrained somewhat by coupling the Fisher information matrix to the covariance of the outputs, as in equation (4.6).

When creating simulations during the ABC procedure we can calculate the true sufficient statistic, i.e.

$$x_i^s = \sum_{j=1}^{n_d} (d_{ij}^s)^2 \quad (6.8)$$

where $i$ labels the random initialisation of the simulation and the $j$ labels the data point in the data set $d$. Plotting the exact sufficient statistic against the network output allows us to see how well the correct function is learned by maximising the IMNN, as seen in figure 4. The blue points show the values of exact sufficient statistics of the simulations and scaled values of the network outputs of the same simulations. The network output must be scaled due to the allowed linear scaling of the sufficient statistic. We actually found that network output is approximately

$$\text{network output} \approx \sum_{j=1}^{n_d} (d_{ij}^s)^2 + 58, \quad (6.9)$$

without a linear scaling of the exact sufficient statistic, but with an offset. The black dashed line shows what would be expected if the exact map was learned by the network. We can see that the network output generally follows the sum of the square of the data closely with hints of a slight bend and superficial broadening at larger exact sufficient statistics. The bending is of no concern since any one-to-one function of the sufficient statistic is still a sufficient statistic, and we can see that the network output is clearly a monotonic function of the real summary. The broadening indicates that only an approximate map is learned because the training of the network is incomplete due to lack of diversity within simulations and perhaps a sub-optimal choice of network hyperparameters. With greater variety within the simulations or, likewise, a greater number of simulations, the optimal map could be learned even more precisely. Nevertheless, we can see how minor an effect the broadening of the exact sufficient statistic is by looking at the results in figure 3. The resulting posterior distribution is equivalent to the analytic posterior, which is the real proof that the network has found the correct summary statistic.

Figure 5: Rescaled network output for a given exact sufficient statistic. The blue dots show a scaled value of the network output at the exact sufficient statistic of simulations obtained at a range of $\vartheta$ during ABC. The black dashed line shows the expected value of the network output if the network had learned the map from data to the sufficient statistic perfectly. Since the scatter of the exact sufficient statistic to the network output closely follows the black dashed line, we know the network has approximately learned the correct map from data to sufficient statistic. There is a slight curve which arises from the fact that any one-to-one function of the sufficient statistic is still a sufficient statistic and so is of no concern. There is also a superficial broadening of the curve which shows that the map is only approximately correct.

1. Summarising Gaussian signals with known noise variance

Now consider some noisy data where the real data $d = \{d_i \sim \mathcal{N}(0, \vartheta + \sigma^2_{\text{noise}}) \mid i \in [1, n_d]\}$ has a signal variance of $\vartheta^{true} = 1$ and the variance of the noise is taken to be known $\sigma^2_{\text{noise}} = 1$. Simulations of the noisy data can be created and used to train the network, as before. The addition of the noise makes the likelihood less peaked about the true parameter value and so the Fisher information is expected to be less than in original problem. Since the likelihood is known analytically, using equation (6.6) it can be seen that $F = 1.25$. The network manages to achieve $F \approx 1.25$ by the end of training, suggesting the network is capable of extracting close to the maximum amount of information possible. We have used a slightly less complex network here with [128, 128], but all other parameters the same. Again, many different architectures work equally well, but do not necessarily
converge as quickly. We train the network for 2000 epochs before the Fisher information saturates to its maximum value.

In figure 6 it can be seen that the PMC-ABC posterior distribution, shown in the purple histogram with shaded 1-σ Poisson regions, when given some simulated test data, \( d \), is very similar to the analytic result shown by the solid orange line. The dashed blue approximate posterior distribution from the asymptotic likelihood again peaks very close to the maximum of the analytic posterior. The posterior distribution becomes maximal at the most likely parameter value given the data, with the variance given by the inverse Fisher information at the end of training. There are 1000 samples used to create the histogram of the PMC-ABC posterior which required approximately \( 2 \times 10^5 \) simulations to be created during the PMC, where all samples are within \( x_{ik}^T = x \pm \varepsilon^T \), with \( \varepsilon^T = 0.109 \). Since the analytic posterior distribution is so similar to the PMC-ABC posterior we can see that, even though the network is only given noisy simulations, it is capable of finding the true function to summarise the data.

2. **Summarising Gaussian signals with unknown noise variance**

Now consider the problem where, again, the real data \( d = \{ d_i \sim N(0, \vartheta + \sigma^2_{\text{noise}}) \mid i \in [1, n_d] \} \) has a signal variance of \( \vartheta = 1 \) and the variance of the noise is also unknown with a uniform prior \( \sigma^2_{\text{noise}} \in (0, 2] \). We train using 1000 simulations (+100 for each of the derivatives) at a fiducial \( \vartheta^{\text{fid}} = 1 \) each with a different \( \sigma^2_{\text{noise}} \) randomly drawn from the uniform prior on the noise. The final value of the Fisher information from the network is less than in either of the two previous cases at \( F = 0.9 \) using a slightly more complex network than in the previous section with an architecture of \([128, 128, 64]\) but all other parameters the same. If the noise were assumed to be known at \( \sigma^2_{\text{noise}} = 2 \) then the maximum Fisher available, as calculated from equation (6.6), would be \( F = 5/9 \). The posterior distributions for \( \vartheta \) are shown in figure 7. Since the noise is unknown, a Rao-Blackwell estimate of the analytic distribution is made. Here, the posterior distribution is calculated for a range of given noise values from \( \sigma^2_{\text{noise}} = (0, 2] \) and their results summed at each value of \( \vartheta \), plotted with a solid orange line. The PMC-ABC posterior is given by the purple histogram consisting of 1000 samples, which required approximately \( 10^5 \) simulations using the PMC. Again, as before, the constraints on \( \vartheta \) are incredibly similar to the analytic result, confirming that the network can approximate the exact summary very well. The Rao-Blackwell estimation procedure is also carried out to obtain the posterior calculated from the asymptotic likelihood, in dashed blue, although the result does not agree with the exact or
PMC-ABC posteriors. The lack of agreement arises because the simulated test data is not well represented in the training simulations. Even though there is an under representation in the data, the network has learned the correct way to summarise data independent of the input, i.e. the network calculates the sum of the square of the input.

B. Summarising Gaussian signals with wrong fiducial variance

Since the network trained in the known noise problem, in section VI.A is akin to a network trained at a fiducial parameter \( \vartheta^{\text{fid}} = 2 \), we can use it to test how well the network can predict the variance when the fiducial value does not coincide with the true parameter. It would be expected that data with \( \vartheta = 1 \) would be under-represented in a training data set where the fiducial value is \( \vartheta^{\text{fid}} = 2 \). Naively, one would assume that the network would not perform as well as a network trained using simulations created at \( \vartheta^{\text{fid}} = 1 \), especially since the Fisher information available from this network is \( F = 1.25 \) and not \( F = 5 \) as in section VI.A. However, figure 8 shows that the parameter constraints given the same real data as in table II are equally as strong as when using the trained network from section VI.A. It is promising that the training of the network seems fairly insensitive to the choice of fiducial parameter. The posterior distribution from the asymptotic likelihood, in dashed blue, is much wider than the same curve in figure 3 since the variance of the distribution is given by the Cramér-Rao bound, i.e. \( F^{-1} = 0.8 \), rather than \( F^{-1} = 0.2 \) when the network from section VI.A. The fact that the purple histogram matches the analytic solid orange distribution so well indicates that the network has learned the correct way to summarise data, rather than learning an algorithm for mapping simulations to an output which specifically depends on the fiducial parameter value. For example, in the problem considered here, we know that the correct summary of the data is the sum of the square of the data (or at least a linear scaling of the sum of the square of the data). The network is trained in such a way that the abstract function of weights, biases and inputs that the network represents closely approximates the sum of the square of the input. Once abstract function is learned, it does not matter what parameter value is used to create the simulations, even if that parameter is far from the fiducial value, because the network will still output the sum of the square of the input. It is extremely encouraging to see that the network can extrapolate beyond its training data by depending on the robustness of the learned patterns.

C. Summarising quasar spectra

Beyond the elementary test case on variance estimation, we can consider models that are of more astronomical interest. Here we attempt to generate constraints on the amplitude of scalar perturbations, \( A_s \), using a simplified 1D model of the Lyman-\( \alpha \) forest from a single quasar. To generate simulations we begin by using the halo mass function calculator hmf module (Murray et al. 2013) in python to generate the 3D power spectrum \( P^\text{lin}(k) \), evolved using the method of Eisenstein and Hu (Eisenstein and Hu 1998), at a redshift of \( z = 2.25 \) with fixed cosmological parameters (at \( z = 0 \)). The cosmological parameters come from the Planck 2015 temperature and low-\( \ell \) polarisation results (Ade et al. 2016), \( H_0 = 67.7 \text{ km Mpc}^{-1}\text{s}^{-1} \), \( \Omega_m = 0.307 \), \( \Omega_b = 0.0486 \), \( n_s = 0.9667 \), \( \sigma_8 = 0.8159 \), \( T_{\text{CMB}} = 2.725 \text{ K} \), \( N_{\text{eff}} = 3.05 \), and \( \sum m_{\nu} = 0.06 \text{ eV} \), calculated using astropy (Astropy Collaboration et al. 2013). The power spectrum is calculated between \( \ln k_{\text{min}}/(1h\text{Mpc}^{-1}) = -18.42 \) and \( \ln k_{\text{max}}/(1h\text{Mpc}^{-1}) = 9.90 \) in steps of \( \Delta \ln k/(1h\text{Mpc}^{-1}) = 0.005 \). The correlation function can be found using

\[
\xi(r) = \int_{0}^{\infty} \frac{dk}{2\pi^2} \exp \left[ -R^2_\eta k^2 \right] k^2 P^\text{3D}(k) \text{sinc}(kr) \tag{6.10}
\]

where the exponential term is a smoothing function where we use \( R_\eta = 5h^{-1}\text{Mpc} \). We calculate the value
of $\xi(r)$ between $-200 < r < 200 h^{-1}\text{Mpc}$ in $N = 8192$ bins. To simulate the density fluctuations along the line of sight, we calculate the 1D power spectrum using

$$P^{1D}(k) = \int_{-\infty}^{\infty} dr \exp[ikr]\xi(r).$$  \hfill (6.11)

The Lyman-$\alpha$ peak in the rest frame of an emitter is $\lambda_{RF}^\text{Ly} = 121.567\text{nm}$ (Bautista et al. 2015) and we use the fact that BOSS can measure absorbers in the redshift range $1.96 < z < 3.44$ (Bautista et al. 2017). Using

$$z = \frac{\lambda}{\lambda_{RF}} - 1$$  \hfill (6.12)

the minimum observed wavelength of the Lyman-$\alpha$ peak is $\lambda_{\text{min}} = 359.838\text{nm}$ (at $z = 1.96$) and the maximum wavelength is $\lambda_{\text{max}} = 539.757\text{nm}$ (at $z = 3.44$) (Bautista et al. 2017). The length $L$ of the survey in comoving space is calculated between these redshifts, yielding $L = 1122.9 h^{-1}\text{Mpc}$. The frequency spacing is given by the inverse of the survey length, so we consider a range of $k = (0, 14.6) h^{-1}\text{Mpc}^{-1}$ with $N = 8192$ bins. We modify the 1D power spectrum such that it more closely follows the gas power spectrum as seen from Lyman-$\alpha$ absorptions (McDonald 2003),

$$P^{1D}_{g}(k) = \beta D(k, \mu)P^{1D}(k)$$  \hfill (6.13)

where $\beta$ is a free parameter, set for a given realisation of noise which ensures that $\langle F \rangle = 0.8$ (Font-Ribera et al. 2012). $D(k, \mu)$ is a term which modifies the small-scale power spectrum (McDonald 2003) and is of the form

$$D(k, \mu) = \exp \left[ \left( \frac{k}{k_{\text{NL}}} \right)^{\alpha_{NL}} - \left( \frac{k}{k_{p}} \right)^{\alpha_{p}} - \left( \frac{k}{k_{V}} \right)^{\alpha_{V}} \right],$$  \hfill (6.14)

where

$$k_{V} = k_{V0} \left( 1 + \frac{k}{k_{p}} \right)^{\alpha_{V}}$$  \hfill (6.15)

and $k_{\text{NL}} = 6.40 h^{-1}\text{Mpc}^{-1}$, $\alpha_{NL} = 0.569$, $k_{p} = 15.3 h^{-1}\text{Mpc}^{-1}$, $\alpha_{p} = 2.01$, $k_{V0} = 1.220$, $k_{V}' = 0.923 h^{-1}\text{Mpc}^{-1}$, $\alpha_{V}' = 0.451$, $\alpha_{V} = 1.50$ (Blomqvist et al. 2015) and we choose to use $\mu = k_{\parallel}/k = 1$ since we only consider independent quasar lines, i.e. the flux is completely decorrelated from one line to the next. The above numbers are computed for the log-flux explicitly described in (McDonald 2003). For the purpose of demonstration we keep the same $k$ dependence here. With the gas power spectrum in equation (6.13), normalised by the length of the survey, we can generate 1D random Gaussian fields, $\delta_g$. The Gaussian fields are generated by multiplying unit variance, zero mean Gaussian noise with $(P^{1D}_{g}(k)/2)^{1/2}$ and Fourier transforming into real space, including the normalisation of $N/(2L)$ due to the discrete nature and finite period of the discrete Fourier transform. The flux from quasars is absorbed by neutral hydrogen in over-densities in the density field, and can be calculated from the fluctuating Gunn-Peterson approximation (Peeples et al. 2010) as

$$F = \exp[-\tau]$$  \hfill (6.16)

where we consider the form of the optical depth to be

$$\tau = 1.54 \left( \frac{T_0}{10^4 \text{K}} \right)^{-0.7} \frac{10^{-12} \text{s}^{-1}}{\Gamma_{\text{UV}}} \left( 1 + z \right)^6 \frac{0.7}{h} \left( \frac{\Omega_b h^2}{0.02156} \right)^2 \frac{4.0972}{H(z)/H_0} \rho^{2-0.7}$$  \hfill (6.17)

where $T_0 = 18400\text{K}$ is the normalisation to the power-law temperature-density relation $T = T_0(1+\delta)^{\gamma}$ with $\gamma = 0.29$ (both here and in equation (6.18)) and $\Gamma_{\text{UV}} = 4 \times 10^{-12}\text{s}^{-1}$ is the photo-ionisation rate due to the ambient UV background (Peeples et al. 2010). The gas density field is normalised such that its mean is unity,

$$\rho = \frac{\exp[\delta_g]}{\langle \exp[\delta_g] \rangle}.$$  \hfill (6.19)

The continuum flux can be calculated between the Lyman-$\alpha$ and Lyman-$\beta$ peaks at $\lambda_{RF}^\text{Ly} = 121.567\text{nm}$ and $\lambda_{RF}^\text{Ly} = 102.572\text{nm}$ using the PCA formulation of (Suzuki et al. 2005). The continuum flux in the rest frame of the emitter is calculated using

$$r(\lambda) = \mu(\lambda) + \sum_i c_i(\lambda)\xi_i(\lambda)$$  \hfill (6.20)

where $\mu(\lambda)$ is the mean flux over many quasars, $\xi_i(\lambda)$ are the $i$th principal components and $c_i(\lambda)$ are the amplitudes of the principal components which we consider to be $c_i(\lambda) = 1$ for simplicity. The continuum can be transformed into the observer’s wavelength space by assuming a redshift for the quasar and inverting equation (6.12). We choose the redshift of the simulated (and real) quasar to be $z = 2.91$. The flux, which is currently in real space, is transformed into wavelength space by interpolating the comoving distance, $r$, along given redshift values, $z$, using the $\text{hmf\ comoving\ distance}(z)$ function and then using equation (6.12). The continuum modulated flux from the quasar is simply

$$f(\lambda) = F(\lambda)C(\lambda)$$  \hfill (6.21)

where $C(\lambda)$ is $r(\lambda)$ from equation (6.20) in the rest frame of the observer (Bautista et al. 2015). Figure 9 shows the generated flux from a single mock quasar at $z = 2.91$ in blue. The orange (lighter) line shows the continuum flux between the Lyman-$\alpha$ and Lyman-$\beta$ peaks and the dashed green line shows the mean of the transmitted flux. We only consider the flux between $406\text{nm} < \lambda < 469.2\text{nm}$ which is $104\text{nm} < \lambda < 120\text{nm}$ in the rest frame of the quasar (Bautista et al. 2017). We bin the wavelengths using the resolution from the BOSS
coadded spectra of $\Delta \log_{10} \lambda/{\rm nm} = 10^{-4}$ (Bautista et al. 2015) which gives a flux in W$m^{-2}$nm$^{-1}$, but needs to be measured in photon counts. Using the method$^2$ in (Ahn et al. 2012) we see that for a quasar such as the one we are generating the spectra for, there is an almost one-to-one correspondence between flux and photon count (albeit the photon count is integer) (Ahn et al. 2012). Therefore, we make the assumption that making the flux into integer values and then applying Poisson noise satisfac-

torily represents real quasar spectra. Our binned, noisy spectra have 581 data points, each of which can be used as an input to an IMNN.

An example of the simulated test data input to the network is shown in figure [10].

For any set of fixed cosmological parameters, $A_s$, is a scaling of $P_s^{1D}(k)$. To get constraints on $A_s$ we can train a network at a fiducial $A_s$ and then use the PMC to find the posterior distribution of $A_s$ compared to some simulated test data. In fact, for simplicity, we can consider the parameter $\vartheta$ to be some multiplicative scaling of the amplitude, $A_s = \vartheta A_{\cosmo}$ with $A_{\cosmo}$ the amplitude of the power spectrum found in equation (6.13). We use $\vartheta^{\text{fid}} = \exp[0]$ as the fiducial parameter, i.e. $A_s^{\text{fid}} = A_{\cosmo}$.

A relatively simple network, such as [256, 256], is able to obtain a Fisher information of $F = 0.015$, which is the maximum Fisher information that could be found over a large range of different network architectures and hyperparameters. However, the network which was most resilient to incorrect fiducial values was more complex than those networks previously considered. The network with the largest Fisher information by the final epoch of training, which could handle incorrect fiducial parameters was a network four hidden layers shaped like $[1024, 512, 256, 128]$, using 1000 simulations (with 100 simulations each for the upper and lower components of the derivative) which were split into two batches, an initial bias of $b = 0.1$, where the activation function is leaky ReLu with $\alpha = 0.1$, a dropout of 20\% and a learning rate of $\eta = 1 \times 10^2$ when training for 10000 epochs.

As before, once the network was trained, PMC-ABC could be performed. We used a uniform prior in logarithmic space of $\vartheta = \exp[-10, 10]$. The simulated test data was created away from the fiducial parameter value of $\vartheta^{\text{fid}} = \exp[0]$ at $\vartheta^{\text{real}} = \exp[3]$, i.e. $A_s = \exp[3]A_{\cosmo}$, and is shown in figure [10]. The posterior distribution for the value of $\vartheta$ can be found in figure [11]. Here, we required 1000 samples in the posterior requiring at least 2500 draws in the final iteration of the PMC to be convinced that the posterior had converged. The histogram peak, and the tentative peak of the leading order expansion of the likelihood, are at their maximum at $A_s \approx \exp[3]A_{\cosmo}$, i.e. $\ln \vartheta \approx 3$, which confirms that the correct test parameter can be recovered, shown as the vertical black dashed line in figure [11]. There is a large, degenerate tail in the PMC-ABC posterior which arises due to the amplitude of the random Gaussian noise, used to create the quasar spectrum, being so small that the features in the generated flux become negligible. Since the network output of the random fluctuations are still reasonably close to the network output from the real data, they cannot be constrained. The lack of constraining power at low $\vartheta$ is even clearer in the posterior from the leading order expansion. The constraints on $A_s$ span approximately 5 orders of magnitude or more using the PMC-ABC posterior, which seems poor, but is due to

2. In particular we use the method described in [http://www.sdss.org/dr12/algorithms/spectrophotometry/](http://www.sdss.org/dr12/algorithms/spectrophotometry/) in the section called “DR9 Flux to Photons”. We use quasar 023918.47+025035.6 as a guideline.
using only one quasar spectrum to constrain cosmology with. Joint inference using several quasars would provide a much stronger constraint, as is done when using cosmological surveys. Although the constraints are not particularly strong, we have shown that we can learn to extract information from highly noisy data, and summarise it in such a way that we can perform PMC-ABC to get a posterior distribution for parameters of interest.

\[ D. \text{ Gravitational waveform frequency} \]

\cite{Graff et al. 2011} showed that the MOPED algorithm, described in section \[ ] was unable to summarise the central oscillation frequency of a gravitational waveform from LISA without introducing spurious features \cite{Graff et al. 2011}. By using non-linear summaries of the data, the problem in \cite{Graff et al. 2011} can be avoided.

We start by considering a sine-Gaussian gravitational wave signal as could be seen in LISA, with a short burst duration and frequency space waveform \cite{Feroz et al. 2010} of

\[
\Lambda(f) = \frac{AQ}{f} \exp \left[ -\frac{Q^2}{2} \left( \frac{f - f_c}{f_c} \right)^2 \right] \exp [2\pi if t_c],
\]

(6.22)

where \( A \) is some amplitude, \( Q \) is the width of the gravitational wave burst, \( t_c \) is the time of the burst and \( f_c \) is the central oscillation frequency. We fix \( A = 3.5, Q = 5 \) and \( t_c = 1 \times 10^5 \) and require that the signal-to-noise of the burst is \( S/N = 34 \) \cite{Graff et al. 2011}. We are interested in summarising and constraining the parameter \( f_c \).

To generate a simulation of the gravitational wave signal, we use the one-sided noise power spectral density of the LISA detector \cite{Feroz et al. 2010}, which is

\[
S_h(f) = 16 \sin^2 \left( 2\pi f t_L \right) \left( 2S_{pn} \left( 1 + \cos \left( 2\pi f t_L \right) \right) + \cos^2 \left( 2\pi f t_L \right) \right),
\]

(6.23)

\[
S_{pn}(f) = \left( 1 + \left( \frac{10^{-4} \text{Hz}}{f} \right)^2 \right) S_{acc},
\]

(6.24)

where \( S_{sn} = 1.8 \times 10^{-37} \text{Hz}^{-1} \) is the shot noise, \( S_{acc} = 2.5 \times 10^{-48} \text{Hz}^{-1} \) is the proof acceleration mass and \( t_L = 16.67 \) is the light travel time along one arm of the LISA constellation. To generate the real space gravitational wave burst, we calculate the frequency space waveform \( \Lambda(f) \) and detector noise \( \pi(f) \) and then Fourier transform them into real space

\[
\Lambda(f) = \int_{-\infty}^{\infty} dt \, \Lambda(t) \exp[2\pi i f t],
\]

(6.25)

\[
\pi(f) = \int_{-\infty}^{\infty} dt \, n(t) \exp[2\pi i f t].
\]

(6.26)

We perform the Fourier transform at 2048 time steps from \( t = 9.9 \times 10^5 \) sampled at 1s intervals. The output of the LISA detector is then given by

\[
d = h(\theta^{true}) + n
\]

(6.27)

where \( h(\theta^{true}) \) is the values of the gravitational waveform at the true parameter values, \( \theta^{true} = \{ A^{true}, Q^{true}, t_c^{true}, f_c^{true} \} \) at the sampled time and \( n \) is a random realisation of the noise. When assuming a noise covariance which is independent of the signal, \( \sigma_n^2 = \mathbb{I} \), the logarithm of the likelihood is particularly simple \cite{Feroz et al. 2010} and is given by

\[
\ln \mathcal{L} = C - \frac{||d - h(\theta)||^2}{2},
\]

(6.28)

where \( h(\theta) \) is the real space gravitational wave at, not-necessarily-true, parameters \( \theta \) and \( C \) is a constant which we set to zero.

We are interested in summarising the data to constrain the central oscillation frequency, \( f_c \), of the gravitational wave. To do so, we use a network which takes in the 2048 inputs from the data with the architecture \{10, 10, 10, 10, 10\}. The network has a 10% dropout and ReLU activation with \( \alpha = 0.01 \). The learning rate is fixed at \( \eta = 10^{-5} \) and the biases are initialised slightly positively at \( b^l = 0.1 \). We train for 1200 epochs using
At different parameters is therefore given by
\[
\ln \mathcal{L} = C - \frac{||x - x^h(f_c)||^2}{2\sigma_n^2}.
\] (6.30)

We calculate equations (6.28) and (6.30) using simulated test data, \(d\), generated at \(f_c^{\text{true}} = 0.1\text{Hz}\) between 1 \times 10^{-2} < f_c < 0.5\text{Hz}. The logarithm of the likelihood of \(f_c\) calculated using all the data, ln \(\mathcal{L}(f_c|d)\) is shown in the upper subplot of figure 12 as a dashed orange line. This is compared to ln \(\mathcal{L}(f_c|x_{MOPED})\) using the MOPED summary as the dotted green line in the middle subplot and ln \(\mathcal{L}(f_c|x)\) using the network summary as the solid blue line in the bottom subplot. The MOPED summary assumes a noise covariance which is independent of the signal such that the compression parameter is simply
\[
r_{f_c} \propto \mu_f f_c.
\] (6.31)

For the network summary, the noise is automatically included through the random initialisation of the simulations used to train the network. It can be seen that each of the likelihoods in figure 12 agree with \(f_c^{\text{true}} = 0.1\text{Hz}\), shown with the dashed black line, but a false aliasing peak appears, shown with the dotted black line, when using the MOPED summary. This false maximum in the likelihood arises from unsuccessfully undoing the Fourier transform which leaves the mapping from \(h(f_c) \rightarrow x^h_{\text{MOPED}}(f_c)\) not being one-to-one. On the other hand, the IMNN compression does not suffer this problem. There is a clear unique summary which, when used to calculate the approximate likelihood assuming Gaussian noise and evaluated at different \(f_c\), results in a single peak at the \(f_c^{\text{true}}\). Full inference on \(f_c\) is then possible using PMC-ABC.

This test shows that, through the use of the non-linear function provided by the IMNN, we are able to surpass the capability of linear compression. Not only can the summary from the network be at least as informative as the MOPED summary, it is also more robust since it is able to avoid misleading parameter inference due to non-trivial mappings.

VII. CONCLUSIONS

We have shown how information maximising neural networks (IMNNs) can perform automatic physical inference. Automatic physical inference begins by training a neural network to find the optimal non-linear summaries of data supplied only with simulations and no other knowledge about how to best compress data. Once the network is trained, its output is used to perform PMC-ABC and find the approximate posterior distribution of any parameter that the network is sensitive to. We have also shown that the network is insensitive to poor choice in fiducial parameter value when generating simulations.

We consider the technique presented in this paper as an extension or replacement to other massive optimal data
compression procedures. The MOPED algorithm is able to optimally compress data using linear combinations under the assumption that the likelihood is known and is, to first order, Gaussian. Further, the method in (Alsing and Wandelt 2017) generalises MOPED to any given likelihood function, where the compressed statistics no longer need to be linear. In (Alsing et al. 2018), the likelihood does not need to be known at all, firstly summarising simulations of real data heuristically and then compressing these summaries using an appropriate likelihood in the same way as (Alsing and Wandelt 2017). Although a powerful technique, the first step in (Alsing et al. 2018) can potentially be lossy and the likelihood in the second step should be well known to achieve optimal compression of the first step summaries. The information maximising neural network can replace both steps in (Alsing et al 2018) by taking the raw data and providing non-linear, likelihood-free summaries directly from the simulations. Likewise, and perhaps more conveniently, the network introduced here is ideally placed to squeeze additional information out of the data after all of the more obvious summaries, such as the power spectrum, have been exhausted.

In this paper, we have focussed on a few test models used to illustrate the method and its abilities. The first set of tests use the network to find a summary of Gaussian signal, without noise, with known noise variance and with unknown noise variance. This is a useful example since it can be solved analytically and linear compression, such as MOPED would fail to provide useful summaries of the data. We showed that PMC-ABC is able to recover the analytic posterior distribution for the variance of the Gaussian noise nearly exactly, which means that the network has correctly learned the sufficient statistic for this problem. It is useful to consider variance inference as there are many examples in astronomy and cosmology where the variance is informative about the underlying parameters. Although the details of the input data and simulations will be more complex, variance estimation appears in cases such as estimating the value of the optical depth to reionisation, $\tau$, and recovering B-mode polarisation from probes of the large-angle cosmic microwave background polarisation anisotropies.

Following the success of the first set of tests, the next two examples show further tests on astronomically motivated problems. The first shows how extremely noisy raw data can be directly input to the network to constrain cosmological parameters and the second shows how using non-linear summaries are suited to situations where linear summaries can be misleading.

Information maximising neural networks are designed to deal with raw data. We can see IMNNs being useful, or even essential, when trying to calculate posterior distributions of model parameters where the likelihood, describing the distribution of some large number of data points, is unknown. For example, the raw data from large scale structure surveys is infeasibly large. Even the number of summary statistics is $\sim 10^4$ and a likelihood cannot be written to describe the physics, the selection bias and the instrument—but the data can in principle be simulated from initial conditions. The IMNNs presented in this paper to illustrate and explore the concept used a fully connected architecture. When considering very large data sets we will need to consider network architectures that are adapted to the problem at hand and computationally efficient. For example, assuming the isotropy of the universe transverse to the line of sight, whilst looking radially in redshift space suggests that stacks of convolutional neural networks could be used to deal with raw LSS data. As long as patches of the large scale structure (and the instrument) can be simulated to train the convolutional filter, IMNNs should make it possible to extract cosmologically interesting information directly from the raw data—automatically.

The data and original code used in this paper is available at 
https://doi.org/10.5281/zenodo.1175196
For up-to-date code and current development please use
https://github.com/tomcharnock/information_maximiser

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