Small and Lightweight Laser Retro-Reflector Arrays for Lunar Landers

Xiaoli Sun\textsuperscript{1,*}, David E. Smith\textsuperscript{2}, Evan D. Hoffman\textsuperscript{1}, Shane W. Wake\textsuperscript{1}, Daniel R. Cremons\textsuperscript{3}, Erwan Mazarico\textsuperscript{1}, Jean-Marie Lauenstein\textsuperscript{1}, Maria T. Zuber\textsuperscript{2}, Edward C. Aaron\textsuperscript{4}

\textsuperscript{1}NASA Goddard Space Flight Center, Code 698/61A/551/561, Greenbelt, MD 20771, USA
\textsuperscript{2}Department of Earth, Atmospheric and Planetary Sciences, Massachusetts Institute of Technology, Cambridge, MA, 02139, USA
\textsuperscript{3}Universities Space Research Association (USRA), Columbia MD 21046, USA
\textsuperscript{4}KBRwyle Technology Solutions, LLC, Lanham, MD 20706, USA

Abstract

A set of small and lightweight laser retro-reflector arrays (LRA) were fabricated and tested for use in lunar landers under NASA’s Commercial Lunar Payload Service (CLPS) program. Each array contains eight 1.27-cm diameter corner cube retro-reflectors mounted on a dome-shaped aluminum structure. The array is 5.0 cm in diameter at the base, 1.6 cm in height, and 21 gram in mass. They can be tracked by an orbiting laser altimeter, such as the Lunar Orbiter Laser Altimeter (LOLA) from a distance of a few hundred kilometers or by a landing lidar on future lunar landers. The LRAs demonstrated a diffraction-limited optical performance. They were tested to survive and function on the Moon for decades, well after the lander missions are completed.

1. Introduction

Small laser retro-reflector arrays (LRA) have been used for laser ranging to Earth-orbiting satellites from ground stations [1]. Large LRAs were placed on the Moon by the Apollo astronauts and the Russian lunar landers and have been used for lunar laser ranging from Earth over the past 50 years [2-4]. Small LRAs have also been considered for use on Mars landers as fiducial markers for future laser tracking from an orbiting lidar [5-7]. Tracking small LRAs on the ground from an orbiting lidar has been demonstrated by the Geoscience Laser Altimeter System (GLAS) on Ice, Cloud, and Elevation Satellite (ICESat) [8]. One LRA by the National Space Agency of Italy (ASI, for Agenzia Spaziale Italiana) was on the ExoMars Schiaparelli module of the European Space Agency (ESA) [9]. A similar LRA from ASI was installed on NASA’s Mars InSight lander for potential future laser ranging from orbit [10]. NASA is planning to put small LRAs on future commercial lunar landers from the United States and international lunar landers for laser ranging from lunar orbiting lidar, such as the Lunar Orbiter Laser Altimeter (LOLA) on board the Lunar Reconnaissance

\textsuperscript{*}xiaoli.sun-1@nasa.gov.
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Orbiter (LRO), which is currently in orbit around the Moon [11, 12]. A set of small LRAs have been fabricated based on the original design for the Phoenix Mars Lander. The arrays are 5.0 cm in diameter, 1.6 cm in height, and weigh 21 gram. They have been tested to withstand the launch and landing and to remain functional on the lunar surface for decades. One of these LRAs was installed on the Beresheet lunar lander by SpaceIL of Israel. Another one has been installed on the Vikram lunar lander on the Chandrayaan-2 mission by the India Space Research Organisation (ISRO) which was launched on July 22, 2019. In this paper, we describe the design and testing of these LRAs and the potential applications for science and exploration.

A small LRA on a lander on the Moon or an asteroid can be detected by an orbiting lidar. The LRA position on the lunar surface can be derived from a series of range measurements from the lidar. Once the location of the LRA has been established, subsequent orbiting lidar can range to the LRA to assist in the orbit determination of the host spacecraft (Figure 1). Over time, these laser ranging measurements can be used to monitor the geological movement of the landing site. The LRA position can also be used to establish a local geodetic coordinate system around the lander.

The small LRAs will remain functional for decades after the lander has accomplished its planned mission objectives. They can continue to be used as precision landmarks for guidance and navigation during the lunar day or night. A landing lidar such as a flash lidar can detect these small arrays during descent. A few LRAs surrounding the landing site can serve as precision land markers to guide the arriving landers. Figure 2 shows how a small number of arrays surrounding a landing pad could be used to help autonomous and safe landing.

Small LRAs can also be deployed as fiducial markers on asteroids, similar to the target markers deployed on the asteroid 162173 Ryugu by the Hayabusa-2 spacecraft [13-15]. One possible configuration is to attach two of these LRAs at the base in a spheroid shape such that one side always face up regardless how the joined arrays land on the asteroid surface. They can then be tracked by an orbiting or co-orbiting spacecraft with a lidar to precisely measure the orbit, rotation speed and axis, and other geophysical properties, as depicted in Fig. 3.

Although these LRAs are small, they have a relatively large equivalent optical cross section of $10^4$ to $10^5$ m$^2$ and can be tracked by a lidar from a few hundred km distance, as described in Section 5. They produce a distinct return which can be uniquely identified from the surface returns by the return pulse amplitude and shape, as shown by ICESat [8]. It may take several orbital passes to be located with today’s topographic lidar, such as LOLA, that has a small laser footprint size and low pulse rate. They are expected to be spotted much easier by future swath mapping lidar with contiguous or overlapping laser footprints in a push-broom measurement configuration.

It should be pointed out that these small LRAs are designed for laser ranging from a lunar orbiting lidar to the lunar lander but not from Earth directly. The retro-reflectors on these LRAs are about 1/3 the size of those on the Apollo LRAs and they are oriented in different
directions for wide sky-view instead of stronger return signals like the Apollo LRAs. The received signal from these small LRA is ≪1% of the signal from the Apollo LRAs. At present, the highest received signal levels from the Apollo LRAs is 6-7 detected photons per laser pulse at the Apache Point Observatory [16]. Thus, it is not practical with today’s technology to range to these small LRAs directly from Earth-based stations.

There are unique technical challenges for the LRAs on lunar landers compared to those on Earth-orbiting satellites or Mars landers. First, they have to be small and lightweight so that they can be carried to the Moon by all new lunar landers from the NASA CLPS program. Second, they have to remain functional over the entire lunar surface temperature range, 85 to 385 K (−189°C to 111°C) [17], which is a much wider range compared to LRAs on Earth-orbiting satellites and landers on Mars. Third, they are exposed to the sky without shielding and have to survive a much higher dose of space radiation and for a much longer period than those required for most spacecraft components.

2. Laser Retro Reflector Array Design

The LRAs were originally designed in the early 2000’s for a Mars Lander for ranging from an orbiting lidar (also known as a laser altimeter) similar to the Mars Orbiter Laser Altimeter (MOLA) [5, 6, 18]. They were specially designed for low mass and a near full sky-view. Each LRA consisted of eight 1.27-cm diameter solid retro-reflectors (also known as corner cubes or cube-corners) mounted on a dome-shaped aluminum structure, as shown in Fig. 4. Each retroreflector has a useful light incident angle of about ±20° [19-22]. Four retro-reflectors are evenly distributed on a ring 20° from zenith and the other four on a ring 40° from zenith. Therefore, the entire LRA can be ranged from any direction 30° above the horizon.

The retro-reflectors are made of a special type of quartz called Suprasil, which has a high optical homogeneity in all three axes and is proven to be radiation resistant. The index of refraction is about 1.46, varying slightly with the wavelength. To reduce the mass, the retro-reflectors are bonded to the aluminum structure instead of mechanically constrained as in those used on Earth orbiting satellites. Standard chromate-conversion coating (Iridite) finish was used for the aluminum structure for high electrical conductivity to prevent the build up of static electrical charges on the LRA in space. The mass of the LRA is about 21 gram, not including fasteners. The original LRA designed for the Mars lander used plated retro-reflectors with the three exterior facets of the corner cubes mirror-coated with aluminum. The recent LRAs for commercial and international lunar landers use non-plated or total internal reflection (TIR), retro-reflectors. Although TIR retro-reflectors have a slightly narrower light acceptance angle and the retro-reflected light is affected by the polarization of the incident light [23], they have a simpler design, are less susceptible to thermal distortion, and are likely to have a longer lifetime on the Moon. The size of the retro-reflectors was chosen to maintain low mass but still give sufficient cross section for laser ranging from an orbiting satellite. The small size and low mass also helps to minimize thermal gradients across the retro-reflector under solar illumination, which are known to cause degradation in optical performance [24].
3. Optical Performance of the Retro-Reflectors

The key parameters that govern the optical performance of the retro-reflectors are the dihedral angles and the surface flatness. They were measured at the retro-reflector supplier and verified at NASA Goddard Space Flight Center (GSFC) using interferometers. The average dihedral angles were measured to be about 0.3 arcseconds root-mean-squares (rms), which is adequate to maintain the optical losses to within 30% [25]. The surface flatness is within 1/10 the laser wavelength at 633 nm. The dihedral angles and surface flatness of one of the retro-reflectors were measured from 100 K to 375 K in 20 K steps and there were no measurable changes over this temperature range.

The far-field patterns of the retro-reflected beams were characterized at NASA GSFC using a 2.5-m focal length collimator system and a camera system at about 2.1 μrad per pixel resolution. The test setup was calibrated against a gold coated mirror of the clear aperture and positioned normal to the incident beam. A sample far-field pattern from a single retro-reflector is shown in Figure 5. The divergence of the center lobe of the retro-reflected beam was measured to be 97 μrad full-angle. The acceptance angle was about ±25° for this plated retro-reflector.

The LRAs showed a nearly diffraction limited far field pattern at both 532 nm and 1064 nm laser wavelength. The details of the optical tests will be reported in a separate publication. These include the dihedral angles and surface flatness of a sample set of single retro-reflectors, dihedral angles and surface flatness of a single retro-reflector versus temperature, far-field patterns from single retro-reflectors and the entire array versus incident angle, far-field patterns at different laser wavelengths and polarizations, and the retro-reflected laser pulse shapes at different incident angle.

4. Environmental Testing

Several environmental tests were performed on the LRAs to verify that they can withstand the vibrations during launch and landing, the temperature range during cruise and on the lunar surface from day to night, and the lunar radiation environment for several decades. A qualification unit was built using the same procedure and materials for some of the tests to avoid unnecessary stress and wear to the flight units. The following subsections give the details of the vibration, thermal vacuum and radiation testing of the LRAs at NASA GSFC before delivering them to the lunar landers.

4.1. Vibration Tests

Vibration tests were performed according to the NASA Standard and the particular lunar lander request. There were two sets of vibration tests conducted: qualification and acceptance tests. The former were performed on the qualification unit to the full vibration levels, including a random vibration test to 26 g rms, a sine sweep test to 100 Hz, and a shock response test to 2000 g. The acceptance tests were performed on all the flight LRAs to verify workmanship per NASA GSFC General Environmental Verification Standard GSFC-STD-7000. The vibration specification for the qualification unit and the flight units are given in Tables 1 through 4, with the Z axis defined to be normal to the LRA base. The LRAs
passed all the vibration tests. A static load test was performed on the qualification unit after the vibration tests and the bond strength between the corner cube retro-reflectors and the aluminum support structure was found to meet the requirement with margin. There were no resonances found over the entire test frequency range.

4.2. Thermal Vacuum Tests

A thermal cycle test was performed for every flight LRA in vacuum over a temperature range of 85 to 385 K. The LRAs were tested in a liquid nitrogen Dewar with an adapter plate and a set of heaters, as shown in Figure 6. The LRAs were held at the hot and cold temperatures for at least 8 hours. The rate of temperature transition was about 3 K/min. There were a total of five temperature cycles for each LRA.

All LRAs passed the thermal vacuum tests. There were no noticeable changes after the tests, except for a slight fading of the Iridite coating on the aluminum surface. Slight discoloring of Iridite is expected under high temperature and vacuum. It is not expected to significantly affect their corrosion resistance during handling and storage on Earth and their lifetime on the lunar surface [26]. Detailed examination of the Iridite coating under a high power microscope showed no metal damage. The bonding strength between the retro-reflectors and the aluminum structure of the qualification unit was also tested after the temperature cycle test and it met the specification with margin.

4.3. Radiation Damage Tests

Since the LRAs are likely to be mounted on the top deck of the landers, they have to face the full space radiation incident on the lunar surface. Most of the concerns about the radiation damage to these LRAs are related to the bonding material, which has not been tested to such a high radiation level. Another radiation damage mechanism to be considered is the darkening of the glasses used for the retro-reflectors. Our goal is to demonstrate a useful lifetime of several decades so that they can be ranged by future lunar lidar.

The LRAs are exposed to space radiation on one side and shielded by the Moon on the other side. The primary radiation is from solar protons when the LRAs are on the sunlit side of the Moon. Protons cause both ionization damage and displacement damage to the materials. For optics like the retro-reflectors, the major radiation damage results from the proton-induced ionizing dose that causes darkening of the material [27, 28]. Since the LRAs are not shielded, most of the ionizing energy is deposited by low energy protons that are absorbed in the top layer of the material, which ionizing the material. Gamma rays are often used to test the ionizing radiation damage susceptibility of optical materials because they are more accessible and lower cost compared to proton accelerators. For the same amount of dose, gamma rays have been shown to have a greater impact on transmission loss in glass as compared with protons [29, 30], thereby providing a more conservative test.

The ionizing dose on the lunar surface over a 10-year period was calculated at the 95% confidence level as a function of aluminum shielding thickness using the Emission of Solar Protons (ESP) model and both the NOVICE and SHIELDOSE software packages. The results are shown in Figure 7. An additional 10% of the total direct dose is added to account for the reflected particles by the lunar surface onto the lunar lander. The total ionizing dose
averaged over 10 years is expected to be 8.8 Mrad(Si) based on the NOVICE result, which is much higher than what most of the spacecraft components are designed for. The adhesive that holds the retro-reflectors to the support structure is mostly behind the aluminum shell and can be assumed to have an equivalent of 1-mm aluminum shielding. The expected dose for the adhesives is below 0.1 Mrad(Si) over 10 years based on Figure 7.

A gamma-ray radiation test was conducted on the qualification unit of the LRA using a Cobalt-60 source. The adhesive and the assembly procedure of the qualification unit were exactly the same as the flight unit, though the retro-reflectors were made of the similar type of quartz materials but not Suprasil used on the flight units. The total dose of the test was 17.8 Mrad(Si) at a dose rate of about 2 krad(Si)/min, which is equivalent to about a 20-year exposure to the space radiation on the lunar surface. The integrity of the adhesives and the clearness of the retro-reflectors were examined periodically during the test and no change was observed except for a slight hardening of the adhesives. A static load test of the retro-reflectors was performed by applying a static force up to 50 times the required retention strength of the adhesive for each of the eight retro-reflectors on the array. The LRA showed no observable differences before and after the radiation test. A single flight-spare retro-reflector was also tested to 13.7 Mrad(Si) and again there was no change found when comparing to an identical but nonirradiated retro-reflector. Based on these results, the adhesives should survive the radiation on the lunar surface for many decades. The LRAs are predicted to have no degradation in optical performance for two decades and remain functional for more decades to come on the lunar surface.

5. Laser Ranging to the LRAs

Link Equation:
The received signal from a single retro-reflector can be calculated using the link equation, as [1],

\[ n_s = \eta_d \eta_r \frac{\lambda}{4\pi R^2} \frac{E_t G_t(\theta_L, \delta\theta_L) \sigma_x(\theta)}{(4\pi R^2)} A_r, \]

where \( n_s \) is the average number of photoelectrons recorded by the detector, \( \eta_d \) is the detector quantum efficiency, \( \eta_t \) and \( \eta_r \) are the transmitter and receiver optical transmission efficiencies, \( \lambda \) is the laser wavelength, \( h \) is Planck’s constant, \( c \) is the speed of light, \( E_t \) is the transmitted laser pulse energy, \( G_t(\theta_L, \delta\theta_L) \) is the transmitter antenna gain as a function of the laser beam divergence angle \( \theta_L \) and the laser transmitter pointing error \( \delta\theta_t \), \( \sigma_x(\theta) \) is the equivalent cross-section of the retro-reflector as a function of the laser beam incident angle \( \theta_L \), \( A_r \) is the receiver telescope light collecting aperture area, and \( R \) is the distance from the lidar to the LRA.

The antenna gain for a Gaussian laser beam with a divergence angle of \( \theta_L \), half angle at \( 1/e^2 \) points, can be written as

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The peak cross-section for a plated retro-reflector under normal incidence can be approximated as [1]

\[
\sigma_{xp} = \frac{\pi^3 \rho \phi_c^4}{4 \lambda^2}
\]

where \(\rho\) is the reflectivity and \(\phi_c\) is the diameter of corner cube retro-reflector. The reflectivity is typically \(\rho = 78\%\) for an aluminum-plated corner-cube retro-reflector. For TIR retro-reflectors, an elliptical polarization is induced into the beam as the light travels inside the retro-reflector. The phase contributions from all six unique paths through the retro-reflector (two paths for each facet of the cube) interfere with one another, changing the overall amplitude and shape of the return. The peak optical cross section for a TIR retroreflector is 26.4\% that of a perfect reflector of the same size [23].

The far-field diffraction pattern (FFDP) of the retro-reflected beam in terms of the light intensity versus the retroreflected angle can be approximated by the Airy function,

\[
\sigma_x(\theta) = \sigma_{xp} \cdot \left( \frac{J_1 \left( \frac{\pi D_c}{\lambda} \theta \right)}{\frac{\pi D_c}{\lambda} \theta} \right)^2
\]

where \(J_1(x)\) is the Bessel function of the first kind of order one. The first null point (half-angle) is given by

\[
\theta_0 = 1.22 \frac{\lambda}{\phi_c}
\]

For the corner cubes we used \((\phi_c = 1.27 \text{ cm})\), the center lobe size is about 50 and 100 \(\mu\text{rad}\) in radius at 0.532 and 1.064 \(\mu\text{m}\) laser wavelengths, respectively.

When the light incidence angle is not normal to the face of the corner cube, the cross section decreases with the angle of incidence as [31]

\[
\sigma_x(\theta_i) = \sigma_{xp} \left( \frac{2 \cos \theta}{\pi} \right)^2 \left[ \sin^{-1} \left( \frac{1 - 2 \tan^2 \left( \frac{\sin \theta_i}{n} \right)}{\sin \left( \frac{\sin \theta_i}{n} \right)} \right) \right]^{-2}
\]
where \( n \) is the index of refraction of the material used for the corner cube. The cross section falls to 10% of the peak value at about \( \pm 25^\circ \).

For a TIR corner cube, the cross section depends not only the incidence angle but also on the polarization and the orientation of the corner-cube [23]. The cross section versus the incident angle has to be calculated numerically [19] [20]. The cross section of TIR retro-reflectors generally have a smaller acceptance angle compared to plated retro-reflectors. The useful incidence angle is about \( \pm 20^\circ \), but can be up to \( \pm 35^\circ \) at certain orientation.

**Velocity Aberration:**

Another major factor to consider is the velocity aberration. Because the lidar on the spacecraft is moving but the return beam is retro-reflected to where the laser pulse was transmitted, the spacecraft is no longer at the peak of the retro-reflected beam by the time the retro-reflected laser pulses return. As a result, the received signal becomes lower than that given by Eq.(1) through (6). The angles of the retro-reflected beam can be calculated as [32]

\[
\theta_r = \acos \left( \frac{(1 - \alpha^2)\cos(\theta_{sc})}{1 + \alpha^2 + 2\alpha \sin(\theta_{sc})\cos(\phi_{sc})} \right),
\]

\[
\phi_r = \atan \left( \frac{(1 - \alpha^2)\sin(\theta_{sc})\cos(\phi_{sc})}{\alpha + \sin(\theta_{sc})\cos(\phi_{sc})} \right),
\]

and

\[
\alpha = \frac{2v_{sc}}{c}.
\]

where \( \theta_r \) and \( \phi_r \) are the angles of the reflected beam, \( \theta_{sc} \) and \( \phi_{sc} \) are the angles of the transmitted laser beam, all in polar coordinate system in reference to the spacecraft position, and \( v_{sc} \) is the spacecraft velocity. The differences between the incident and the reflected beam angles, \( \Delta \theta_r = \theta_r - \theta_{sc} \) and \( \Delta \phi_r = \phi_r - \phi_{sc} \) gives the angular deviation from the peak of the retro-reflected beam. The parameter \( \alpha \) given in Eq.(8) is often used as a close approximation to these angular deviations.

For the 1.27-cm diameter retro-reflector at \( 10^\circ \) in the along-track and cross-track directions (i.e. \( \theta_{sc} = \phi_{sc} = 10^\circ \)), and a typical spacecraft velocity in lunar orbit (1.6 km/s), the angular deviations are \( \Delta \theta_r = 9.71 \mu \text{rad} \) and \( \Delta \phi_r = 10.0 \mu \text{rad} \), which is about 1/10 the half angle of the retro-reflected beam divergence at 1.064 \( \mu \text{m} \) laser wavelength. Therefore, the effect of velocity aberration is usually not a major factor for laser ranging to these small laser retro-reflectors from a lidar in a typical lunar orbit.

The maximum ranging distance from the spacecraft to these laser retro-reflector arrays on the lunar landers can be estimated using the above equations once the instrument parameter values are known. As an example, we consider using LOLA to range to these retro-reflectors with the relevant parameters listed in Table 5 [11]. For a \( \theta_{sc} = \phi_{sc} = 10^\circ \) and assuming only one retro-reflector is contributing to the return signal, the maximum ranging distance is 300
km when LOLA was operating at its full laser pulse energy. The LOLA laser pulse energy has been decreasing over time and the far-field beam pattern has been degrading to multiple spatial modes [12, 33]. The current LOLA laser output pulse energy is estimated to be 1/10 of the original value and the maximum ranging distance to these laser retro-reflectors is about 170 km. The major challenge for LOLA to range to these LRAs is the targeting of laser beam to the LRAs. Since LOLA was designed for surface topography measurements, the laser beam divergence angle was designed for high spatial resolution instead of wide coverage. The laser pulse rate is 28 Hz due to the constraints on the electrical power. LOLA samples the lunar surface with five 5-m diameter laser spots at 25-m spacing from a 50-km orbit altitude and could easily miss the laser retro-reflectors on any given pass. Furthermore, the laser beam pointing is controlled by the spacecraft body pointing, which has a relatively large uncertainty in real time during the maneuver and could miss the LRAs on a given pass. Therefore repeated attempts are required in order for LOLA to range to these small LRAs on the lunar surface. These difficulties can be overcome by future lunar lidar with higher sensitivity and pixelated detectors and contiguous ground coverage.

6. Summary

A set of small and lightweight laser retro-reflector arrays (LRA) were fabricated and tested for use in lunar landers under NASA’s Commercial Lunar Payload Service (CLPS). Each LRA contains eight 1.27-cm diameter corner cubes on a dome shaped aluminum structure. The LRA is 5.0-cm in diameter at the base, 1.6 cm in height and 21 gram in mass. The arrays were tested to survive the vibration during launch and landing. They were tested over the extreme temperature range on the lunar surface, from of 85 K to 385 K. They were tested with gamma rays to a total dose of 17.8 Mrad(Si) and showed no measurable damage. The LRAs demonstrated a near diffraction limited far-field patterns from the optical performance tests. These LRAs can be tracked by a LOLA-like orbiting laser altimeter from a few hundred km distance, or by a lander lidar on the lunar lander during descent. These LRAs on the lunar landers are expected to remain functional and serve as the fiducial marks on the Moon for many decades to come.

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Fig. 1.
Laser ranging from an orbital lidar to determine the lander position on the lunar surface, and help the orbiter to self-determine its trajectory once the position of lander is established.
Fig. 2.
Laser ranging to an array of LRAs as precision landmarks for safe landing by returning landers.
Fig. 3. Deployable LRAs on an asteroid for laser ranging from an orbiting or co-orbiting spacecraft with a lidar to determine its orientation, rotation axis, etc.
Fig. 4.
Photograph of the LRA. It is 5.0 cm in diameter at the base and 1.6 cm in height, and weighs 21 g.
Fig. 5.
Far field pattern of a plated (mirror-coated) retro-reflector at 532 nm laser wavelength.
Fig. 6.
LRA thermal vacuum test setup.
Fig. 7.
Average total ionization dose on lunar surface over 10 years.
Table 1.

Random vibration test level - qualification

| Frequency Range | Level       | Units  |
|-----------------|-------------|--------|
| All 3 axes      |             |        |
| 20 Hz           | 0.04 g^2/Hz |        |
| 80-300 Hz       | 1.5 g^2/Hz  |        |
| 300-1300 Hz     | 0.3 g^2/Hz  |        |
| 1300-3000 Hz    | 1.5 g^2/Hz  |        |
| Overall level   | 26 g rms    |        |
| Duration        | 2 minutes per axis |
Table 2.

Sine vibration test level - qualification

| Frequency   | Z axis | X&Y axes |
|-------------|--------|----------|
| 5 - 20 Hz   | 8.3 g  | 8.3 g    |
| 20 - 50 Hz  | 13.3 g | 13.3 g   |
| 50 - 100 Hz | 60 g   | 30 g     |
| Sweep rate  | 3-4 octaves/minute | 3-4 octaves/minute |
Table 3.

Shock response test - qualification

| Frequency | Z axis | X&Y axes |
|-----------|--------|----------|
| 100 Hz    | 30 g   | 30 g     |
| 1,500 Hz  | 2,000 g| 2,000 g  |
| 10,000 Hz | 2,000 g| 2,000 g  |
## Table 4.

Random vibration test level - workmanship

| Frequency Range | Test Level | Spectral Slope |
|-----------------|------------|----------------|
| All 3 axes      | 0.013 g²/Hz| +6 dB/octave   |
| 20 Hz           | 0.013 g²/Hz| +6 dB/octave   |
| 20-50 Hz        | 0.08 g²/Hz | +6 dB/octave   |
| 50-800 Hz       | 0.08 g²/Hz | +6 dB/octave   |
| 800-2000 Hz     | 0.08 g²/Hz | +6 dB/octave   |
| 2000 Hz         | 0.013 g²/Hz|                |
| Overall level   | 10 g rms   |                |
| Duration        | 1 minute per axis |    |
**LOLA Instrument Parameters**

| Parameters                  | Symbols | Values               |
|-----------------------------|---------|----------------------|
| Laser pulse energy          | $E/\text{beam}$ | 3.3 mJ into 5 beams |
| Divergence angle            | $\theta_L$ | 100 μrad             |
| Wavelength                  | $\lambda$ | 1.064 μm             |
| Detector quantum efficiency | $\eta_d$ | 40%                  |
| Transmitter optical transmission | $\eta_t$ | 95%                  |
| Receiver telescope diameter | $\phi_{tel}$ | 0.14 m             |
| Receiver optical transmission | $\eta_r$ | 70%                  |
| Laser pointing error        | $\delta \theta_t$ | 15 μrad             |
| Retro-reflector diameter    | $\phi_c$ | 1.27 cm              |
| Minimum detectable signal   | - | 90 photoelectrons/pulse |
| Retro-reflector optical efficiency | $\rho$ | 25%                  |
| Spacecraft velocity         | $\nu_{sc}$ | 1.6 km/s            |