Silicon spin qubits have emerged as a promising path to large-scale quantum processors. In this prospect, the development of scalable qubit readout schemes involving a minimal device overhead is a compelling step. Here we report the implementation of gate-coupled rf reflectometry for the dispersive readout of a fully functional spin qubit device. We use a p-type double-gate transistor made using industry-standard silicon technology. The first gate confines a hole quantum dot encoding the spin qubit, the second one a helper dot enabling readout. The qubit state is measured through the phase response of a lumped-element resonator to spin-selective interdot tunneling. The demonstrated qubit readout scheme requires no coupling to a Fermi reservoir, thereby offering a compact and potentially scalable solution whose operation may be extended above 1 K.
The recent years have witnessed remarkable progress in the development of semiconductor spin qubits\textsuperscript{1–4} with an increasing focus on silicon-based realizations\textsuperscript{5–8}. Access to isotopically enriched\textsuperscript{28}Si has enabled the achievement of very long spin coherence times for both nuclear and electron spins\textsuperscript{4–11}. In addition, two-qubit gates with increasing high fidelities were demonstrated in electrostatically defined electron double quantum dots\textsuperscript{12–14}. While further improvements in single- and two-qubit gates can be expected, growing research efforts are now being directed to the realization of scalable arrays of coupled qubits\textsuperscript{15–19}. Leveraging the well-established silicon technology may enable facing the scalability challenge, and initiatives to explore this opportunity are on the way\textsuperscript{20}. Simultaneously, suitable qubit device geometries need to be developed. One of the compelling problems is to engineer scalable readout schemes. The present work addresses this important issue.

It has been shown that a microwave excitation applied to a gate electrode drives Rabi oscillations via the electric-dipole spin resonance mechanism\textsuperscript{4–6,21–23}. The possibility of using a gate as a sensor for qubit readout would allow for a compact device layout, with a clear advantage for scalability. Gate reflectometry probes charge tunneling transitions in a quantum dot system through the dispersive shift of a radiofrequency (rf) resonator connected to a gate electrode\textsuperscript{24–27}. Jointly to spin-selective tunneling, e.g. due to Pauli spin blockade in a double quantum dot (DQD), this technique provides a way to measure spin states.

In a similar fashion, the phase shift of a superconducting microwave resonator coupled to the source of an InAs nanowire has enabled spin qubit dispersive readout\textsuperscript{22}. In Si, recent gate reflectometry experiments have shown single-shot electron spin detection\textsuperscript{28–30}.

Here, we combine coherent spin control and gate dispersive readout in a compact qubit device. Two gates tune an isolated hole DQD, and two distinct electric rf tones (one per gate) allow spin manipulation and dispersive readout. Spin initialization and control are performed without involving any charge reservoir; qubit readout relies on the spin-dependent phase response at the DQD charge degeneracy point. We assess hole single spin dynamics and show coherent spin control, validating a protocol for complete qubit characterization exploitable in more complex architectures.

Results

Double quantum dot dispersive spectroscopy. The experiment is carried out on a double-gate, p-type Si transistor fabricated on a silicon-on-insulator 300-mm wafer using an industry-standard fabrication line\textsuperscript{6}. The device, nominally identical to the one in Fig. 1c, has two parallel top gates, G\textsubscript{R} and G\textsubscript{C}, wrapping an etched Si nanowire channel. The gates are defined by e-beam lithography and have enlarged overlapping spacers to avoid doping implantation in the channel. The measurement circuit is shown in Fig. 1a. At low temperature (we operate the device at 20 mK using a dilution refrigerator), DC voltages V\textsubscript{C} and V\textsubscript{R} are applied to these gates to induce two closely spaced hole quantum dots. The control gate G\textsubscript{C} delivers also sub-\textmu s pulses and microwave excitation in the GHz range to manipulate the qubit. The readout gate, G\textsubscript{RE}, is wire-bonded to a 220 nH surface-mount inductor. Along with a parasitic capacitance and the device impedance, the inductor forms a tank circuit resonating at f\textsubscript{0} = 339 MHz. Figure 1b shows the phase \( \phi \) and attenuation \( A \) of the reflected signal as a function of the resonator driving frequency \( f\textsubscript{b} \). From the slope of the phase trace at \( f\textsubscript{0} \) we extract a quality factor \( Q\textsubscript{loaded} \approx 18 \). The qubit device acts as a variable impedance load for the resonator, and the resonant frequency \( f\textsubscript{0} \) undergoes a dispersive shift according to the state of the qubit.

To determine the charge stability diagram of our DQD, we probe the phase response of the resonator while sweeping the DC gate voltages \( V\textsubscript{R} \) and \( V\textsubscript{C} \) (see Supplementary Note 2 and Supplementary Fig. 2).

The diagonal ridge in Fig. 2a denotes the interdot charge transition we shall focus on hereafter. Along this ridge, the electrochemical potentials of the two dots line up enabling the shunting of a hole charge from one dot to the other. This results in a phase variation \( \Delta \phi \) in the reflected signal. Quantitatively, \( \Delta \phi \) is proportional to the quantum capacitance associated with the gate voltage dependence of the energy levels involved in the interdot charge transition. Interdot tunnel coupling results in the formation of molecular bonding (+) and anti-bonding (−) states with energy levels \( E\textsubscript{+} \) and \( E\textsubscript{−} \), respectively. These states have opposite quantum capacitance since \( C\textsubscript{Q,\pm} = -\alpha^2(\partial^2 E\textsubscript{\pm}/\partial V^2)^{27} \). Here \( \epsilon \) is the gate-voltage detuning along a given line crossing the interdot charge transition boundary, and \( \alpha \) is a lever-arm parameter relating \( \epsilon \) to the energy difference between the electrochemical potentials of the two dots (we estimate \( \alpha \approx 0.58 \text{ eV V}^{-1} \) along the detuning line in Fig. 2a). The width of the...
\( \Delta \phi \) ridge, once translated into energy, gives the interdot tunnel coupling, \( t \). We estimate \( t \) between 6.4 and 8.5 \( \mu \)eV, depending on whether thermal fluctuations contribute or not to the dispersive response (see Supplementary Note 3).

The total charge parity and the spin character of the DQD states can be determined from the evolution of the interdot ridge in an applied magnetic field, \( B \). Figure 2b shows the \( B \)-dependence of the phase signal at the detuning line indicated in Fig. 2a. Four representative traces taken from this plot are shown in Fig. 2c. The interdot phase signal progressively drops with \( B \), as revealed by the trace at \( B = 0.85 \) T. The observed behavior can be understood in terms of an interdot charge transition with an even number of holes in the DQD, in a scenario equivalent to a \((0, 2) \leftrightarrow (1, 1)\) transition. We shall then refer to a “\((0, 2)\)” and a “\((1, 1)\)” state, even if the actual number of confined holes is larger (see Supplementary Note 2). The \( \epsilon \) dependence of the DQD states at finite \( B \) is presented in Fig. 2d. Deeply in the positive detuning \( \epsilon \), the ground state is a spin-singlet state \( S(0, 2) \) and the triplet states \( T(0, 2) \) lie high up in energy. Around zero detuning, the \( | ↑ \uparrow \rangle, | ↓ \uparrow \rangle \) states hybridize with the \( S(0, 2) \) state forming an unpolarized triplet \( T_0(1, 1) \) and two molecular singlets, \( S_0 \) and \( S_r \), with bonding and anti-bonding character, respectively (Supplementary Note 3).

We use the spectrum of Fig. 2d to model the evolution of the interdot phase signal in Fig. 2b, c. Importantly, we make the assumption that the average occupation probability of the available excited states are populated according to a Boltzmann distribution with an effective temperature \( T_{\text{eff}} \) which is used as a free parameter. Because the reflectometry signal is averaged over many resonator cycles, \( \Delta \phi = \sum_i \langle \Delta \phi \rangle_i \), where \( \langle \Delta \phi \rangle_i \) is the phase response associated to state \( i \) weighted by the respective occupation probability (Supplementary Note 3) (here \( i \) labels the DQD levels in Fig. 2d). Figure 2e shows \( \langle \Delta \phi \rangle_i \) as a function of \( \epsilon \) for \( T_{\text{eff}} = 250 \) mK. The spin polarized triplet states \( T_− \) and \( T_+ \) (i.e. \( | ↑ \uparrow \rangle \) and \( | ↑ \downarrow \rangle \), respectively) are linear in \( \epsilon \) and, therefore, they do not cause any finite phase shift; \( S_\sigma \) and \( T_0(1, 1) \), on the other hand, possess a curvature and are sensed by the reflectometry apparatus (Supplementary Note 3). We note that the phase signal for \( T_0(1, 1) \) has a peak-dip line shape whose minimum lies at positive \( \epsilon \) (blue trace), partly counterbalanced by the positive phase signal due to \( S_\sigma \). The \( S_\sigma \) state causes a pronounced dip at negative \( \epsilon \) (green trace), dominating over the peak component of \( T_\| \). The overall net result is a phase signal with an asymmetric double-peak structure consistent with our experimental observation.

This simple model, with the chosen \( T_{\text{eff}} = 250 \) mK, qualitatively reproduces the emergence of the double-dip structure at \( B = 0.4 \) T, as well as its gradual suppression at higher \( B \), as shown in the inset to Fig. 2b and in Fig. 2f (increasing the Zeeman energy results in the depopulation of the \( S_\sigma \) and \( T_0(1, 1) \) excited states in favor of the \( T_−(1, 1) \) ground state, for which \( \Delta \phi = 0 \)).

**Dispersive detection of electric-dipole spin resonance.** Now that we have elucidated the energy-level structure of the DQD, we can discuss the operation of the device as a single-hole spin qubit with electrical control and dispersive readout. Electric-dipole spin resonance (EDSR) is induced by a microwave voltage modulation applied to gate \( G_C \). To detect EDSR dispersively, the resonating states must have different quantum capacitances.
We ascribe the resonance line to a second-harmonic driving process where $2hf_C = gh_0B$ ($h$ the Planck’s constant, $\mu_B$ the Bohr magneton and $g$ the effective hole $g$-factor). From this resonance condition we extract $g = 1.735 \pm 0.002$, in agreement with previous works. The first harmonic signal, shown in the inset to Fig. 3a, is unexpectedly weaker. Though both first and second harmonic excitations can be expected, the first harmonic EDSR line (inset to Fig. 3a) is unexpectedly weaker. A comparison of the two signal intensities requires the knowledge of many parameters (relaxation rate, microwave power, field amplitude) and calls for deeper investigations.

The visibility of the EDSR signal can be optimized by a fine tuning of the gate voltages. Figure 3b shows a high-resolution measurement over a narrow region of the stability diagram along the interdot charge transition boundary at $B = 0.52$ T; the interdot line has a double-peak structure, consistently with the data in Fig. 2b, c. The measurement is performed while applying a continuous microwave tone $f_C = 7.42$ GHz. EDSR appears as a distinct phase signal around $V_C \approx 362.5$ mV and $V_R \approx 1040$ mV, i.e. slightly inside the $(1,1)$ charge region, pinpointed by the black arrow as I/R. Such EDSR feature is extremely localized in the stability diagram reflecting the gate-voltage dependence of the hole $g$-factor.

Figure 3c displays line cuts across the interdot transition line at fixed $V_R$ and different microwave excitation conditions. With no microwaves excitation, we observe the double-peak line shape discussed above. With a microwave gate modulation at $f_C = 7.42$ GHz, the spin resonance condition is met at $V_C \approx 362.45$ mV, which results in a pronounced EDSR peak, the same observed at point I/R in Fig. 3b (see also Supplementary Fig. 4). The peak vanishes when $f_C$ is detuned by 20 MHz (cyan trace).

At point I/R, resonant microwave excitation enables the spectroscopy of the $T_0(1,1)$ state. The inset to Fig. 3c shows the signal we expect from our model (Supplementary Note 4). In a small detuning window, the populations of $T_0(1,1)$ and $T_0(1,1)$ are assumed to be balanced by EDSR (see the energy levels in the inset to Fig. 3b); this results in a phase signal dramatically enhanced resembling the feature centered at I/R in the main panel. A further confirmation that the spin transitions are driven between $T_0(1,1)$ and $T_0(1,1)$ is given by the extrapolated intercept at 0 T of the EDSR transition line in Fig. 3a, found much smaller (<100 MHz) than $T_0$. In the following, we shall use point I/R to perform qubit initialization and readout.

**Qubit control and readout.** The device is operated as a spin qubit implementing the protocol outlined in Fig. 4a. The voltage sequence in the upper part of Fig. 4a tunes the DQD at the control point C (≈1 mV deep in the $(1,1)$ region) where holes are strongly localized in either one or the other dot with negligible tunnel coupling. A microwave burst of duration $\tau_{\text{burst}}$ and frequency $f_C$ drives single spin rotations between $|\uparrow\downarrow\rangle$ and $|\uparrow\uparrow\rangle$; the system is then brought back to I/R in the “shallow” (1, 1) regime for a time $t_{\text{wait}}$ for readout and initialization. The dispersive readout eventually relies on the spin-resolved phase shift at I/R, though the reflectometry tone $f_b$ is applied during the whole sequence period $T_M$ and the reflected signal is streamed constantly to the acquisition module.

First, we determine the lifetime $T_1$ of the excited spin state at the readout point I/R by sweeping $t_{\text{wait}}$ after a $\pi$-burst at point C. The results are shown in Fig. 4b. The phase signal rapidly diminishes with increasing $t_{\text{wait}}$ because spin relaxation depopulates the excited spin state in favor of the non-dispersive $T_0(1,1)$ ground state. The estimated spin lifetime at the readout position is $T_1 = 2.7 \pm 0.7$ µs (see Supplementary Note 5). By shifting the position of a 100 ns microwave burst within a 12 µs pulse, no
clear decay of the dispersive signal is observed, which suggests a spin lifetime at manipulation point longer than 10 μs. We demonstrate coherent single spin control in the chevron plot of Fig. 4c. The phase signal is collected as a function of microwave burst time τburst and driving frequency fC. The spin state is initialized at point I/R (τwait~T1). In Fig. 4d the phase signal is plotted as a function of τburst set at the Larmor frequency fLarmor. The Rabi oscillations have 15 MHz frequency, consistent with Refs. roro and gtensor. The non-monotonous envelope is attributed to random phase accumulation in the environment. As recently shown42, gate-controlled DQDs can serve as spin qubits even if placed at environmental temperatures exceeding the spin splitting, like 1 K or more. This should relax many cryogenic constraints and support the co-integration with classical electronics, as required by a scale-up perspective19.

Discussion

The measured T1 is compatible with the relaxation times obtained for hole singlet-triplet states in acceptor pairs in Si36 and in Ge/Si nanowire double quantum dots37; in both cases T1 has been measured at the charge degeneracy point with reflectometry set-ups similar to ours. Nonetheless, charge detector measurements have shown T1 approaching 100 μs for single hole spins in Ge hut wire quantum dots38 and ≤1 ms for Ge/Si singlet-triplet systems39. This suggests that despite the intrinsic spin-orbit coupling single spin lifetimes in the ms range might be achievable in Si too.

Strategies to boost T1 at the readout point may consist of inserting rf isolators between the coupler and the amplifier to reduce the backaction on the qubit and avoiding high-κ dielectric in the gate stack to limit charge noise.

We note that T1 could depend on the orientation of the magnetic field as well40. Future studies on magnetic field anisotropy will clarify whether T1 along with the effective g-factors (and hence the dispersive shift for readout) and Rabi frequency, can be maximized at once along a specific direction. Technical improvements intended to enhance the phase sensitivity, like resonators with higher Q-factor and parametric amplification, could push the implemented readout protocol to distinguish spin states with a micro-second integration time, enabling single shot measurement as reported in a recent experiment with a gate-connected superconducting resonant circuit41. Lastly, the resonator integration in the back-end of the industrial chip could offer the possibility to engineer the resonant network at a wafer scale, guaranteeing controlled and reproducible qubit-resonator coupling.

The gate-based dispersive sensing demonstrated here does not involve local reservoirs of charges or embedded charge detectors. This meets the requirements of forefront qubit architectures (e.g. Ref. 16), where the spin readout would be performed at will by any gate of the 2D quantum dot array by frequency multiplexing.

Experimental setup. The experiment is performed by exciting the resonator input at fL = fR = 339 MHz and power P0 = −110 dBm. We measure the phase variation ∆ϕ of the reflected signal isolated from the incoming wave by a directional coupler, amplified by 35 dB at 4 K and demodulated to baseband using homodyne detection. The complete circuit diagram of the experimental setup for qubit manipulation and dispersive readout is provided in Supplementary Note 1 and Supplementary Fig. 1.

Data availability

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Received: 24 December 2018 Accepted: 22 May 2019

NATURE COMMUNICATIONS | https://doi.org/10.1038/s41467-019-10848-z ARTICLE

Fig. 4 Single spin control and dispersive sensing. a The pulse sequence alternating between “deep” (1, 1) regime (C) for spin manipulation and “shallow” (1, 1) regime (I/R), close to the (0, 2) → (1, 1) transition, for the readout and resetting of the spin system. A microwave burst rotates the hole spin during the manipulation stage. The readout tone is continuously applied throughout the sequence period Twait. b Phase shift as a function of τwait for a ±1 mV pulse on Vc with τburst = 100 ns and fC = 12.865 GHz, with B = 0.512 T along β = 0° and θ = 60°. The phase signal approaches 0 when τburst ≫ T1. A simple model yields T1 = 2.7 ± 0.7 μs. c Dispersive signal Δϕ (fC, τburst), measured with the detuning pulses of a) with τwait = 1 μs. Four maps have been averaged. d Phase response as a function of EDSR burst time at fC = 12.865 GHz. The plot shows Rabi oscillations with 15 MHz frequency due to coherent spin rotations. Each data point is integrated for 100 ms and then averaged over 30 traces.

References

1. Petta, J. R. et al. Coherent manipulation of coupled electron spins in semiconductor quantum dots. Science 309, 2180–2184 (2005).
2. Laird, E. A., Pei, F. & Kouwenhoven, L. A valley–spin qubit in a carbon nanotube. Nat. Nanotechn. 8, 565–568 (2013).
3. Dutt, M. G. et al. Quantum register based on individual electronic and nuclear spin qubits in diamond. Science 316, 1312–1316 (2007).
4. Watzinger, H. et al. A germanium hole spin qubit. Nat. Commun. 9, 3902 (2018).
5. Kawakami, E. et al. Electrical control of a long-lived spin qubit in a si/sige nanowire double quantum dot. Nat. Nanotechn. 9, 666–670 (2014).
6. Maurant, R. et al. A CMOS silicon spin qubit. Nat. Comm. 7, 13575 (2016).
7. Reed, M. et al. Reduced sensitivity to charge noise in semiconductor spin qubits via symmetric operation. Phys. Rev. Lett. 116, 110402 (2016).
8. Thorgrimsson, B. et al. Extending the coherence of a quantum dot hybrid qubit. npj Quant. Inform. 3, 32 (2017).

Methods

Device fabrication. The fabrication process of the device was carried out in a 300 mm CMOS platform and is described in Ref. 6.

Experimental set-up. The experiment is performed by exciting the resonator input at fL = fR = 339 MHz and power P0 = −110 dBm. We measure the phase variation ∆ϕ of the reflected signal isolated from the incoming wave by a directional coupler, amplified by 35 dB at 4 K and demodulated to baseband using homodyne detection. The complete circuit diagram of the experimental setup for qubit manipulation and dispersive readout is provided in Supplementary Note 1 and Supplementary Fig. 1.
9. Veldhorst, M. et al. An addressable quantum dot qubit with fault-tolerant control-fidelity. Nat. Nanotechn. 9, 380–383 (2012).
10. Crippa, A. et al. Electrical spin driving by g-matrix modulation in spin-orbit qubits. Phys. Rev. Lett. 120, 137702 (2018).
11. Cacarelli, C. & Ferguson, A. Impedance of the single-electron transistor at radio-frequencies. New J. Phys. 13, 093015 (2011).
12. Collens, J. et al. Dispersive readout of a few-electron double quantum dot with fast rf gate sensors. Phys. Rev. Lett. 110, 046805 (2013).
13. Gonzalez-Zalba, M., Barraud, S., Ferguson, A. & Betz, A. Probing the limits of gate-based charge sensing. Nat. Commun. 6, 6084 (2015).
14. Crippa, A. et al. Level spectrum and charge relaxation in a silicon double quantum dot probed by dual-gate reflectometry. Nano. Lett. 17, 1001–1006 (2017).
15. Paskiam, P. et al. Single-shot single-gate rf spin readout in silicon. Phys. Rev. X 8, 041032 (2018).
16. West, A. et al. Gate-based single-shot readout of spins in silicon. Nat. Nanotechn. 14, 437–441 (2019).
17. Urdampilleta, M. et al. Gate-based high fidelity spin read-out in a CMOS device. Preprint at https://arxiv.org/abs/1809.04584 (2018).
18. Schroer, M., Jung, M., Petersson, K. & Petta, J. Radio frequency charge parity meter. Phys. Rev. Lett. 109, 166802 (2012).
19. Nadj-Perge, S., Frolov, S., Bakkers, E. & Kouwenhoven, L. P. Spin-orbit qubit in a semiconductor nanowire. Nature 468, 1084–1087 (2010).
20. Nadj-Perge, S. et al. Spectroscopy of spin-orbit quantum bits in indium antimonide nanowires. Phys. Rev. Lett. 108, 166801 (2012).
21. Venutucci, B., Bourdet, L., Pouzada, D. & Niquet, Y.-M. Electrical manipulation of semiconductor spin qubits within the g-matrix formalism. Phys. Rev. B 98, 155319 (2018).
22. Scarfino, P. et al. Second-harmonic coherent driving of a spin qubit in a sil/sige quantum dot. Phys. Rev. Lett. 115, 106802 (2015).