ON THE ORIGIN OF DYNAMICALLY COLD RINGS AROUND THE MILKY WAY

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ABSTRACT

We present a scenario for the production of dynamically cold rings around the Milky Way via a high-eccentricity, flyby encounter. These initial conditions are more cosmologically motivated than those considered in previous works. We find that the encounters we examine generically produce a series of nearly dynamically cold ringlike features on low-eccentricity orbits that persist over timescales of \( \sim 2-4 \) Gyr via the tidal response of the primary galaxy to the close passage of the satellite. Moreover, they are both qualitatively and quantitatively similar to the distribution, kinematics, and stellar population of the Monoceros Ring. Therefore, we find that a high-eccentricity flyby by a satellite galaxy represents a cosmologically appealing scenario for forming kinematically distinct ringlike features around the Milky Way.

Subject headings: galaxies: interactions — galaxies: structure — Galaxy: kinematics and dynamics — Galaxy: structure — methods: n-body simulations

Online material: color figures

1. INTRODUCTION

We live in a hierarchical universe, in which mergers are a frequent occurrence (Lacey & Cole 1993; Somerville & Kolatt 1999; Somerville et al. 2000). In this scenario, galaxies like the Milky Way (MW) build up much of their mass by accreting smaller, satellite galaxies. These minor mergers are connected to a variety of observable stellar structural and kinematic phenomena in the MW and external galaxies, including the buildup of stellar halos (e.g., Bullock & Johnston 2005; Bell et al. 2007), “antitruncated” (Erwin et al. 2005; Pohlen & Trujillo 2006; Younger et al. 2007) and “extended” outer disks (Ibata et al. 2005; Peñarrubia et al. 2006; Ibata et al. 2007), and the dynamical heating of the stellar disk (e.g., Toth & Ostriker 1992; Quinn et al. 1993; Walker et al. 1996; Kazantzidis et al. 2007).

In our own Galaxy, Newberg et al. (2002) recently identified a coherent ringlike structure in Monoceros (MRi), using data from the Sloan Digital Sky Survey (SDSS; York et al. 2000). Since this initial detection, the MRi has been identified in the infrared from the Two Micron All-Sky Survey (2MASS; Skrutskie et al. 2006) by Rocha-Pinto et al. (2003) and Martin et al. (2004), reanalyzed using updated SDSS data (Grillmair et al. 2006; Belokurov et al. 2006), and followed up photometrically by several optical surveys (Ibata et al. 2003; Conn et al. 2005b, 2007). The MRi subtends \( \approx 100^\circ \) in Galactic longitude and lies at a galactocentric radial distance of between 15 and 20 kpc. Spectroscopic studies have shown that it is kinematically distinct from the disk—dynamically cold with a low-eccentricity orbit—and is composed primarily of low-metallicity stars with \( [\text{Fe/H}] \approx -0.4 \) (Crane et al. 2003; Yanny et al. 2003; Conn et al. 2005a; Martin et al. 2005, 2006).

Numerical modeling has suggested that similar structures could be formed via tidal disruption of a low-mass satellite on a nearly coplanar orbit (Helmi et al. 2003; Meza et al. 2005; Martin et al. 2005; Peñarrubia et al. 2005, 2007). The most successful of these models argues for a very low mass companion with \( M_{bg}/M_{MW} \sim 1/1000 \) on a nearly circular \( (e = 0.1) \) prograde orbit (Martin et al. 2005; Peñarrubia et al. 2005, 2007). However, recent cosmological simulations (Benson 2005; Khochfar & Burkert 2006) have found that the orbits of satellite galaxies are likely to be highly eccentric. The dynamical friction timescale for such interactions is very long, and as a result it is unlikely that such a satellite could have circularized in time to form the MRi (Besla et al. 2007; Boylan-Kolchin et al. 2008). Therefore, while they are successful at reproducing many of the characteristics of the MRi, these models may have cosmologically unappealing initial conditions.

With this in mind, we propose an alternative mechanism for forming dynamically cold ringlike structures around the MW: a flyby encounter with a small satellite on a high-eccentricity orbit. This scenario was first suggested by Kazantzidis et al. (2007), but here we investigate the triggering event in detail and discuss its relevance to the MRi.

2. SIMULATIONS

The simulations presented in this study were performed with Gadget2 (Springel 2005), an N-body/SPH (smooth particle hydrodynamics) code using the entropy-conserving formalism of Springel & Hernquist (2002). We include radiative cooling and star formation, tuned to fit the observed local Schmidt law (Schmidt 1959; Kennicutt 1998). We also incorporate a sub-resolution multiphase model of the interstellar medium (ISM) (Springel & Hernquist 2003)—softened \( (q_{\text{gas}} = 0.25) \) such that the mass-weighted ISM temperature is \( \sim 10^{24} \) K—and sink particles representing supermassive black holes that accrete gas and release isotropic thermal feedback to self-regulate their growth (Springel et al. 2005b).

The progenitor galaxy models were constructed following Springel et al. (2005a), to which we refer the reader for details. The primary galaxy (total mass of \( M_{200} = 10^{12} h^{-1} M_\odot \)) is analogous to the MW with a baryonic mass fraction of \( m_b = 0.05 \). It was realized with \( 10^6 \) halo particles in a Hernquist (1990) profile with a concentration of \( c = 9 \) as motivated by
3. DISCUSSION

The interaction is summarized in Figures 1 and 2. After the initial close passage, resonances between the orbital frequency of the satellite and primary galaxy disk particles ("stars") excite coplanar tidal arms (Toomre & Toomre 1972). These features then wrap around the primary galaxy as it continues to revolve, forming a set of concentric ringlike features. Star formation induced by the merger is concentrated in the nucleus of the primary galaxy (see, e.g., Hernquist 1989; Mihos & Hernquist 1994, 1996; Hernquist & Mihos 1995), making these rings primarily disk stars. They are furthermore dynamically cold and kinematically distinct from typical disk stars (see Fig. 3) and slowly disperse on a timescale of \( \sim 2-4 \) Gyr (or several rotation periods) owing to phase mixing of the collisionless stellar particles (Binney & Tremaine 1987). Because these rings are generated via gravitational interactions, they are largely insensitive to the gas content or structural parameters of the satellite galaxy. And because the flyby is effectively an impulse interaction, they will still be formed even if the satellite is disrupted during the encounter.

4. RELEVANCE TO THE MONOCEROS RING

The ring features that are produced by the flyby interaction are similar in several ways to the MRi first identified by Newberg et al. (2002). They provide a good match to the kinematics and location of the ring and are roughly consistent with metallically measurements of MRi stars. And cosmological simulations indicate that such interactions are common for MW-sized halos over the timescale for ring formation (i.e., within the past \( \sim 4 \) Gyr; Stewart et al. 2007). Therefore, while we do not claim to have modeled the MRi in detail, on the basis of these similarities, we propose a flyby interaction as a possible formation scenario.

The MRi forms a coherent structure over \( \sim 100^\circ \) in Galactic longitude from approximately 15–20 kpc from the Galactic center. The ring is characterized by a sharp inner edge and a smooth outer edge, with a characteristic width of \( \sim 10^3 \) kpc. The ring is associated with a high surface density of stars, particularly in the central portion, indicating a recent burst of star formation. The MRi is also associated with a high concentration of gas, suggesting a recent interaction with a gas-rich companion.

Cosmological simulations (Bullock et al. 2001) and \( 4 \times 10^5 \) stellar disk (80% of the baryonic mass) and \( 2 \times 10^5 \) gas (20% of the baryonic mass) particles. The satellite galaxy had a total mass of \( M_{\text{halo}} = 5 \times 10^{10} h^{-1} M_{\odot} \) \((M_{\text{pc}}/M_{\odot} \sim 20)\) and an identical baryonic mass fraction. It was realized with \( 10^7 \) halo particles in a Hernquist (1990) profile with a concentration of \( c=18 \), again motivated by cosmological simulations, and \( 2 \times 10^4 \) stellar disk (50% of the baryonic mass) and \( 4 \times 10^4 \) gas (50% of the baryonic mass) particles. They were placed on a parabolic encounter with \( R_p = 5h^{-1} \) kpc perigalactic radius (\( \sim 1 \) scale length), consistent with the results of cosmological simulations (Benson 2005; Khochfar & Burkert 2006).

![Fig. 1.—Evolution of the projected stellar mass density, colored according to a logarithmic scale. The panels are 70 kpc on a side, and the simulation time is printed in the upper left-hand corner in units of Gyr. The color bar indicates the projected \( K \)-band surface brightness in units of mag arcsec\(^{-2}\), assuming a constant ratio of 2 in solar units. [See the electronic edition of the Journal for a color version of this figure.]](image1)

![Fig. 2.—Three different projections (X-Y, Z-Y, and X-Z) of the \( K \)-band surface brightness (left; again assuming a constant \( M/L \) ratio of 2 in solar units) and gas surface mass density (right) at \( t \approx 3 \) Gyr. Panels are 50 kpc on a side in the X and Y directions and 20 kpc in the Z direction. The solid line corresponds to one of the ring structures at roughly the same galactocentric location \( (R_{\text{pc}} = 19 \) kpc\) as the MRi. [See the electronic edition of the Journal for a color version of this figure.]](image2)
center (Newberg et al. 2002; Ibata et al. 2003; Rocha-Pinto et al. 2003; Martin et al. 2004; Conn et al. 2005b, 2007). Kinetically, it is distinguished from MW disk stars by its low radial velocity dispersion ($\sim 25$ km s$^{-1}$); the MRi is dynamically cold (Crane et al. 2003; Yanny et al. 2003; Conn et al. 2005a; Martin et al. 2005, 2006). Proper-motion measurements\(^8\) also suggest that the ring stars are in low-eccentricity orbits (e.g., Crane et al. 2003; Yanny et al. 2003). We find that our simulations produce a ring in the right galactocentric radial range that is approximately circularly supported—$0.9 \leq \Omega / Q_p \leq 1.1$, where $\Omega$ is the orbital frequency and $Q_p$ is that expected from circular motion—with a similarly low radial velocity dispersion (see Fig. 3). This ring also extends $\sim 4 \leq z \leq 4$ kpc above and below the Galactic plane, which is consistent with detections such as those of Conn et al. (2007). Although it is unclear what fraction of the ring would be identified in observations, it contains $\lesssim 1\%$ ($\sim 5 \times 10^3 M_\odot$) of the total stellar mass of the disk, which is consistent with the mass estimates of Yanny et al. (2003).

A disk origin for MRi stars is also broadly consistent with their observed stellar population. Several spectroscopic studies (e.g., Yanny et al. 2003) have found mean metallicities of $[\text{Fe/H}] = -1.6 \pm 0.3$, while Crane et al. (2003) find a higher mean metallicity of $[\text{Fe/H}] = -0.4$ using a different tracer population. This suggests a primarily metal-poor stellar population with a spread in metallicities and possibly multiple epochs of star formation. Our modeling indicates that the MRi may have formed from outer-disk stars moved outward by tidal interactions with the satellite galaxy, which is consistent with this metal-poor stellar population (Luck et al. 2006; Yong et al. 2006). Furthermore, the tidal interaction moves a significant supply of cold gas into the same ring structures (see Fig. 2), which provides the raw material for subsequent epochs of star formation. While the density-dependent prescription in our simulations does not accurately capture some modes of star formation owing to the limitations of the SPH method, prescriptions that include an approximate treatment of shocks (e.g., Barnes 2004) may show these multiple star formation episodes.

While ring features in our simulations capture many of the properties of the MRi, there are some observations that potentially conflict with our proposal. In particular, the satellite in our simulations is roughly the same mass as the LMC and by $\sim 2$ Gyr after the interaction would be at a galactocentric distance of $\sim 250$ kpc. Such a massive object is well within the detection threshold of current surveys (e.g., Willman et al. 2002; Koposov et al. 2007); Leo I—a far lower mass MW companion—was detected at comparable distance (Caputo et al. 1999; Bellazzini et al. 2004; Mateo et al. 2007). However, dynamically cold ring structures are produced generically in flyby encounters, and similar features are present both in a lower inclination interaction ($i = 10^\circ$) and for a lower mass satellite galaxy ($M_{1/200} / M_{PG} \sim 1/100$; Fig. 4). It is thus possible that the satellite galaxy that produced the MRi is either (1) hidden from view, at low Galactic latitude, and/or on the other side of the Galaxy, or (2) of lower mass with $100 \lesssim M_{1/200} / M_{PG} \lesssim 20$. Therefore, it is not unreasonable to speculate that Leo I, which had at least one passage through the MW disk within the past $2-4$ Gyr (Sohn et al. 2007), could have excited the MRi we see today. Moreover, it is likely that this interaction stripped a significant fraction of its mass, which makes current dynamical estimates ($M_{1/100} / M_{PG} \sim 1/1000$; Mateo et al. 2007) a lower limit on its mass at the time of the interaction. Furthermore, simulations indicate that there are likely several subhalos in the MW halo with $100 \lesssim M_{100} / M_{PG} \lesssim 1/20$ on orbits that take them through the disk (Moore et al. 1999; Giocoli et al. 2007; Kazantzidis et al. 2007). Since galaxy formation may be inefficient in low-mass halos (e.g., Haiman et al. 1997; Barkana & Loeb 1999; Kravtsov et al. 2004), they may be significantly underluminous and therefore difficult to identify.

Finally, there have been suggestions in the literature (Crane et al. 2003) measure an overall circular velocity of $\sim 100$ km s$^{-1}$. However, we note (as in Peharubia et al. 2005) that the accuracy of such measurements in physical units is limited to $\Delta v_{circ} \approx 4.74 R_\odot \Delta \mu$, where $R_\odot \approx 8$ kpc is the galactocentric radius of the Sun and $\Delta \mu \sim 3-4$ mas yr$^{-1}$ is the typical proper-motion measurement error. We have therefore chosen to interpret these measurements as being generally consistent with low-eccentricity orbits.

\(^8\) In our simulations, both the 1/20 and 1/100 encounters stripped $\sim 50\%$ of the satellite mass. While the efficiency of this process is sensitive to both the initial structure of the satellite and the orbital parameters of the interaction, this suggests that a passage through the disk is likely to remove a significant fraction of the satellite’s total mass.
et al. 2003; Frinchaboy et al. 2004) that some globular clusters (GCs) may be associated with the MRI. This is also potentially in conflict with our modeling. However, we note that their physical association with the MRI is somewhat speculative. Alternatively, there are several nearly coplanar GCs at roughly the same radius (Harris 1996) as progenitor MRI stars in our simulations that they could have also been moved into the rings by the interaction.

5. CONCLUSION

We present a scenario for the production of dynamically cold rings around a disk galaxy such as the MW via a prograde flyby encounter with a satellite galaxy. Tidal arms excited during close passage coalesce and wrap around the disk of the primary galaxy. These kinds of interactions are more cosmologically likely than the nearly circular orbits presented by Martin et al. (2005) and Peñarrubia et al. (2005, 2007) while dynamical friction is insufficient to circularize the orbit of such a low-mass companion. Our modeling shows rings with similar spatial distribution and kinematics to the MRI. The disk origin of MRI stars is furthermore broadly consistent with observed stellar populations. Therefore, we find that a flyby encounter represents a more cosmologically appealing scenario for the production of the MRI and other dynamically cold rings around the MW.

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Fig. 4.—Three projections (X-Y, Z-Y, and X-Z) of the K-band surface brightness (agin assuming a constant M/L ratio of 2 in solar units) at the same simulation time as Figs. 2 and 3 for a interaction with iden-