Nonlinear Magnetoelectric Response of Fe$_{73.5}$Cu$_1$Nb$_3$Si$_{13.5}$B$_9$/Piezofiber Composite for a Pulsed Magnetic Field Sensor

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Abstract: In this paper, we report the nonlinear magnetoelectric response in a homogenous magnetostrictive/piezoelectric laminate material. The proposed magnetoelectric stack Fe$_{73.5}$Cu$_1$Nb$_3$Si$_{13.5}$B$_9$/piezofiber is made up of high-permeability magnetostrictive Fe$_{73.5}$Cu$_1$Nb$_3$Si$_{13.5}$B$_9$ foils and a piezoelectric Pb(Zr, Ti)O$_3$ fiber composite. The time dependence of magnetoelectric interactions in the Fe$_{73.5}$Cu$_1$Nb$_3$Si$_{13.5}$B$_9$/piezofiber structure driven by pulsed magnetic field was investigated in detail. The experimental results show that the magnetoelectric effect is strongly dependent on the external bias magnetic and pulsed magnetic field parameters. To detect the amplitude of a pulsed magnetic field, the output sensitivity reaches 17 mV/Oe, which is excited by a 100 µs width field. In addition, to measure the pulsed width, the output sensitivity reaches 5.4 mV/µs in the range of 0–300 µs. The results show that the proposed Fe$_{73.5}$Cu$_1$Nb$_3$Si$_{13.5}$B$_9$/piezofiber sensor is ideally suited for pulsed magnetic field measurement.

Keywords: magnetoelectric response; pulsed magnetic field sensor; piezofiber; composite

1. Introduction

Magnetic field sensors are widely used in various fields, for example, in the industry, agriculture, medicine, aerospace, marine, exploration, and drilling. In recent years, the magnetoelectric (ME) composites with a product property of magnetostrictive and piezoelectric phases have been of great importance for the realization of a magnetic sensor [1–3]. In the ME effect, an electric field is induced under the application of a magnetic field or conversely, magnetization is induced under the application of an external electric field [1–3]. To date, different magnetic sensors based on ME composites have been experimentally and theoretically investigated in order to obtain the best ME coupling through changing the structures [4–7], materials [8–10], number of layers [11,12], etc.

Among the proposed composites, the 2-2 laminated composite material has a good ME coupling effect at room temperature, a large degree of freedom in design, and a strong application prospect, which provides a clear physical meaning for the design of a new generation of electronic devices [1–12]. Recently, for static or quasi-static magnetic field sensing in an unshielded room temperature and pressure and lab environment, a high direct current (DC) magnetic field sensitivity of 2.8 Hz/nT and a limit of detection of 800 pT were reported in the NEMS AlN/FeGaB resonator [13]. For alternating current (AC) magnetic field sensing, an extremely low equivalent magnetic noise of 5.1 pT/√Hz at 1 Hz was reported in the Metglas/piezofiber structure [14]. A super-high magnetic sensitivity of $1.35 \times 10^{-13}$ T was directly detected at 23.23 kHz in the 1D (1-1) connectivity ME composites of Metglas/PMN-PT [15].
In practice, transient pulsed magnetic field measurement exists in many fields, such as lightning current in power systems [16,17] and aircrafts [18,19], crack detection with pulsed magnetic flux leakage techniques [20,21], and transient electromagnetic measurement in transformer substations [22]. However, to date, most of the reported theoretical and experimental studies carried out on ME sensors have been investigated under the premise of static magnetic field and standard sine wave magnetic field excitation [1–15]. Only a few studies have focused on the transient response of ME materials in pulsed magnetic fields [23]. The magnetic–mechanical–electric nonlinear coupling mechanism of the magnetostrictive/piezoelectric structure should be different by changing the parameters of transient magnetic fields. Thus, for potential applications in pulsed magnetic field measurement of the ME composite, this paper focused on the transient nonlinear ME response under a pulsed magnetic field. The influence of the pulsed magnetic field parameters and external bias field on the ME structure were investigated in detail.

2. Experimental

The ME structure consists of Fe$_{73.5}$Cu$_1$Nb$_2$Si$_{13.5}$B$_9$ (FeCuNbSiB) and piezofiber, as shown in Figure 1a,b. The FeCuNbSiB (International standard trademark 1K107, produced by Foshan Huaxin Microlite Metal Co., Ltd., Foshan, China) is a pizeomagnetic phase with a high permeability ($\mu_r = 30,000$), high saturation magnetization ($\mu_0 M_s = 1.45$ T), and a large anisotropic constant (~30,000 J/m$^3$). The dimensions of the FeCuNbSiB foils are $10 \times 6 \times 0.025$ mm$^3$. The Pb(Zr, Ti)O$_3$ fiber composite (piezofiber) (M4010-P1, provided by Smart Material Cor., Sarasota, FL, U.S.A.) is a piezoelectric phase consisting of rectangular piezoceramic rods (40 mm long and 180 μm thick) sandwiched between layers of adhesive, electrodes, and polyimide film. The electrodes are attached to the film in an interdigitated pattern which transfers the applied voltage directly to and from the ribbon-shaped rods.

As shown in Figure 1a, two FeCuNbSiB layers (each made by three FeCuNbSiB foils) were subsequently laminated to both the top and bottom surfaces of the Pb(Zr, Ti)O$_3$ (PZT)-fiber layer. Under an AC magnetic field, a mechanical strain was generated in the FeCuNbSiB layer and was then transferred to the piezofiber, resulting in the generation of charges due to the direct piezoelectric effect. As shown in Figure 1b, the ME FeCuNbSiB/piezofiber was designed to operate as a half-wavelength longitudinal resonator vibrating freely at both ends. Thus, one mechanical anchor (made by Beryllium bronze plate, produced by Shanghai Dayu Metal Products Co., Ltd., Shanghai, China) bonded with epoxy at the middle of the FeCuNbSiB/piezofiber where the displacement was zero.

![Figure 1](image_url)

**Figure 1.** (a) Schematic illustrations of the FeCuNbSiB/piezofiber magnetoelectric (ME) composites, (b) photograph of the FeCuNbSiB/piezofiber ME composites, (c) schematic illustrations of the experimental setup.
The experimental setup is shown in Figure 1c. A signal generator (Tektronix AFG3021B, Tektronix Inc., Beaverton, OR, USA) provided a controllable input current to a long straight solenoid coil with 1800 turns and a 182 mm length, which was used to provide the pulsed magnetic field. A Helmholtz coil (Linkphysics Co., Ltd., Shanghai, China) was used to provide the DC bias magnetic field, which was driven by a power amplifier (KEPCO Bipolar Operational Power Supply, KEPCO Inc., Flushing, NY, USA.). The ME output voltages of the FeCuNbSiB/piezofiber laminate were measured with a lock-in amplifier (SR-830, SRS, Sunnyvale, CA, USA.) and an oscilloscope (Tektronix Inc., Beaverton, OR, USA). The bias magnetic field \( H_{\text{bias}} \) was measured with a Gauss meter. The AC magnetic field was calculated in the experiments. The resistance of the long straight solenoid coil was measured as 37.8 \( \Omega \) by a multimeter (Fluke Corporation, Everett, WA, USA). In addition, the voltage of the long straight solenoid in the experiments was measured by an oscilloscope. Thus, the actual current \( I \) in the long straight solenoid can be calculated as \( I = \frac{voltage}{37.8} \). Then, the AC magnetic field of the central solenoid \( H = n*I \) (A/m), where \( n = \frac{1800}{0.182} \).

3. Results and Discussion

The ME response of the FeCuNbSiB/piezofiber composite was calibration-tested under a standard sine magnetic field. Figure 2a shows the ME coefficient \( \alpha_{\text{ME}} \) and the phase angle as a function of the bias magnetic field \( H_{\text{bias}} \) driven at a 1 kHz sine magnetic field. As shown in Figure 2a, \( \alpha_{\text{ME}} \) increased with increasing \( H_{\text{bias}} \) up to about \( H_{\text{bias}} = 3.6 \) Oe reached a maximum value of \( \alpha_{\text{ME}} = 9.2 \) mV/Oe and then gradually decreased as \( H_{\text{bias}} \) was further increased. The induced voltage was independent of \( H_{\text{bias}} \) history and no offset value was found near \( H_{\text{bias}} = 0 \) Oe. These characteristics are quite important to magnetic field detection. In addition, as the direction of \( H_{\text{bias}} \) was changed, a 180° phase shift was found, as shown in Figure 2a. Next, \( \alpha_{\text{ME}} \) was measured as a function of the AC magnetic field frequency at \( H_{\text{bias}} = 5.2 \) Oe while sweeping near the mechanical resonance, as shown in Figure 2b. As this figure shows, the fundamental resonant frequency for the FeCuNbSiB/piezofiber composite was \( \sim 26.47 \) kHz. At this resonant frequency, a value of \( \alpha_{\text{ME}} > 50 \) mV/Oe was reached. The natural period of the FeCuNbSiB/piezofiber composite \( T = \frac{1}{f_r} = \sim 37.78 \) \( \mu \)s.

For the ME composite consisting of mechanically coupled magnetostrictive and piezoelectric layers, the resonance frequency of the ME composite is [12]

\[
f_r = \frac{1}{2l} \sqrt{\frac{1}{\rho \mathbf{s}}}
\]

where \( l \) is the length, and \( \rho \) and \( \mathbf{s} \) are the average density and the equivalent elastic compliance, respectively. The length of the PZT-fiber is 40 mm. The \( \rho_p \) and \( s_{11}^E \) of PZT-fiber is 5.44 g/cm\(^3\) and 32.96 \( \times 10^{-12} \) m\(^2\)/N, respectively. The length of the FeCuNbSiB foils is 100 mm. The \( \rho_m \) and \( s_{33,m}^H \) of FeCuNbSiB is 7.25 g/cm\(^3\) and 5.2 \( \times 10^{-12} \) m\(^2\)/N, respectively. The calculated first-order longitudinal resonant frequencies of the PZT-fiber and the FeCuNbSiB foil are \( \sim 29.5 \) kHz and \( \sim 25.8 \) kHz, respectively. Thus, the resonant frequency for the FeCuNbSiB/PZT-fiber stack should be about \( \sim 25.8 \) kHz to \( \sim 29.5 \) kHz. In order to calculate the resonance frequency of FeCuNbSiB/piezofiber, we used the ANSYS 19.0 software (ANSYS, Inc., Canonsburg, PA, USA). In the simulations, the elastic modulus and density of the epoxy layer were 3 GPa and 5 g/cm\(^3\), respectively. The simulation results of the first-order longitudinal vibration are shown in the inset of Figure 2b. The first-order longitudinal resonant frequency of the FeCuNbSiB/PZT-fiber was \( \sim 28.14 \) kHz. The theory calculated and the simulation results agree with the experimental results.
The time self-oscillation of 51 Oe and \( \Delta t \) was excited by the same amplitude pulsed magnetic field of 51 Oe when \( \Delta t \) times were 0–250 \( \mu s \). Moreover, the other interesting result from Figure 4 is that the FeCuNbSiB/piezofiber structure always oscillates at the natural period \( t = T \) s. After \( t > \Delta t \) s, the pulse magnetic was ~8.4 ns. From this figure, the \( V_o \) increased gradually when the pulsed magnetic field increased rapidly to an amplitude of 51 Oe and maintained ~100 \( \mu s \). This is because the magnetic energy increased gradually when the pulsed field decreased rapidly to zero. However, the \( V_o \) decreased gradually and oscillated at the period \( \Delta t_2 = 37.78 \mu s \). This is due to the inertia effect of the FeCuNbSiB/piezofiber structure. The FeCuNbSiB/piezofiber ME structure is a mechanical resonant structure that self-vibrates near the equilibrium position at the natural period \( T \) when the pulsed magnetic field vanishes. It is clear that the FeCuNbSiB/piezofiber structure always oscillates at the natural period \( T \) in the time domain. If the width of the pulsed magnetic field \( \Delta t \) is equal to \( T \) or multiples of \( T \), what will happen?

Figure 4 shows the time dependence \( V_o \) of the FeCuNbSiB/piezofiber structure driven by magnetic field pulses with amplitudes of 51 Oe and widths of \( \Delta t = T - 5T \). The inset of Figure 4 shows the detail when times were 0–250 \( \mu s \). After \( t > \Delta t \), it is interesting that the self-vibration phenomena vanished when \( \Delta t = T - 5T \), which is extremely different from Figure 3. This result can be explained by the fact that the FeCuNbSiB/piezofiber structure was exactly in an equilibrium position when the pulse magnetic field vanished. Moreover, the other interesting result from Figure 4 is that \( V_o \) grew exponentially, which was excited by the same amplitude pulsed magnetic field of 51 Oe when \( \Delta t = T - 5T \). This result demonstrates that the ME FeCuNbSiB/piezofiber structure can be used as a pulse magnetic width measured device, and this will be investigated in detail in the following section.

**Figure 2.** The ME voltage coefficient of the FeCuNbSiB/piezofiber laminates: (a) ME coefficient \( \alpha_{ME} \) as a function of the bias magnetic field \( H_{bias} \) in response to a ~1 kHz sine driving magnetic field, and (b) \( \alpha_{ME} \) as a function of the AC sine magnetic drive frequency sweeping through the electromechanical resonance.

Then, the transient nonlinear ME coupling was investigated in detail. Figure 3 shows the time dependence output voltage \( V_o \) of the FeCuNbSiB/piezofiber structure driven by the magnetic field pulse with a width (\( \Delta t \)) of ~100 \( \mu s \) and an amplitude of 51 Oe. The inset of Figure 3 shows the detail when times were 0–200 \( \mu s \). The rise time of the pulse magnetic was ~8.4 ns. From this figure, the \( V_o \) increased gradually when the pulsed magnetic field increased rapidly to an amplitude of 51 Oe and maintained ~100 \( \mu s \). It is clear that \( V_o \) oscillated at the period \( \Delta t_1 = 37.78 \mu s \). After \( t > \Delta t = 100 \mu s \), the pulse magnetic field decreased rapidly to zero. However, the \( V_o \) decreased gradually and oscillated at the period \( \Delta t_2 = 37.78 \mu s \). This is due to the inertia effect of the FeCuNbSiB/piezofiber structure. The FeCuNbSiB/piezofiber ME structure is a mechanical resonant structure that self-vibrates near the equilibrium position at the natural period \( T \) when the pulsed magnetic field vanishes. It is clear that the FeCuNbSiB/piezofiber structure always oscillates at the natural period \( T \) in the time domain. If the width of the pulsed magnetic field \( \Delta t \) is equal to \( T \) or multiples of \( T \), what will happen?
Figure 3. Time dependence output voltage $V_o$ of the FeCuNbSiB/piezofiber structure driven by the magnetic field pulse with a width ($\Delta t$) of $\sim$100 $\mu$s and an amplitude of 51 Oe. The inset shows the detail when times were 0–200 $\mu$s. The rise time of the pulse magnetic was $\sim$8.4 ns. The bias magnetic field was set to zero.

Figure 4. Time dependence output voltage $V_o$ of the FeCuNbSiB/piezofiber structure driven by the magnetic field pulse with an amplitude of 51 Oe and a width of $\Delta t = T-5T$. The rise time of the pulse magnetic was $\sim$8.4 ns. The bias magnetic field was set to zero.
Based on Figure 2a, the ME response of the FeCuNbSiB/piezofiber structure is dependent on the $H_{bias}$. Therefore, the $H_{bias}$ dependence of the ME response for the FeCuNbSiB/piezofiber structure driven by the pulsed magnetic field was also investigated. Figure 5 shows the peak output voltage $V_p$ (maximum $V_o$ in Figure 3) as a function of $H_{bias}$ for the FeCuNbSiB/piezofiber structure driven by the pulsed magnetic field with $\Delta t = -100$ μs. The curve tendency in this figure is the same as that in Figure 2a. The maximum $V_p$ was ~0.9 V at $H_{bias} = 3.6$ Oe.

![Figure 5](image)

**Figure 5.** The peak output voltage $V_p$ (maximum $V_o$ in Figure 3) as a function of $H_{bias}$ for the FeCuNbSiB/piezofiber structure driven by the pulsed magnetic field with $\Delta t = -100$ μs.

Next, the transient sensing characteristics of the FeCuNbSiB/piezofiber structure was investigated. Firstly, the amplitude of the pulsed magnetic field $H_A$ sensitivity was measured. Consequently, Figure 6 shows the peak output voltage (like the peak in Figure 3) as a function of $H_A$ at $H_{bias} = 0$ Oe. The linear fitting expression based on the experimental data was $V_p = 0.017H_A - 0.0433$. Based on Figure 6, it is obvious that the induced ME voltage had a near linear ($R^2 = 0.9675$) relation with $H_A$. The observation indicates an improved detection sensitivity of 17 mV/Oe. Clearly, the presented FeCuNbSiB/piezofiber sensor seems to be an ideal application for the detection of amplitude variations of pulsed magnetic fields.

From Figure 4, changing $\Delta t$ results in the variation of the output voltage of the FeCuNbSiB/piezofiber structure. Figure 7 shows the peak output voltage (like the peak in Figure 3) as a function of the width of the pulsed magnetic field $\Delta t = 0$–500 μs as $H_{bias} = 0$ Oe. It is evident that the $V_p$ increases with $\Delta t$ and reaches its maximum value after $\Delta t > 300$ μs. For the linear fitting at $\Delta t = 0$–300 μs, the relationship between $V_p$ and $\Delta t$ is $V_p = 0.0054\Delta t + 0.212$. One can see that the linear ($R^2 = 0.9757$) dependence of the $V_p$ on the $\Delta t$ takes place in a region of $0 < \Delta t < 300$ μs. Based on the slope of the plots, the sensitivity of the FeCuNbSiB/piezofiber structure for pulsed width sensing was determined to be $5.4$ mV/μs. It can be concluded that this ME sensor is suitable for pulsed width measurement fields, such as lightning current monitoring.
Figure 6. The peak output voltage $V_p$ (maximum $V_o$ in Figure 3) as a function of the amplitude of the pulsed magnetic field $H_A$ for the FeCuNbSiB/piezofiber structure driven by the pulsed magnetic field with $\Delta t = -100 \, \mu s$. The red line is the linear approximation.

Figure 7. The peak output voltage $V_p$ (maximum $V_o$ in Figure 3) as a function of the width of the pulsed magnetic field $\Delta t$ for the FeCuNbSiB/piezofiber structure at $H_{bias} = 0 \, \text{Oe}$. The blue line is the linear approximation. The amplitude of the pulsed magnetic field is 52.3 Oe.

4. Conclusions

In conclusion, a FeCuNbSiB/piezofiber structure prototype was fabricated and several experiments were conducted to illuminate the transient nonlinear ME response. The time dependence output voltage of the FeCuNbSiB/piezofiber structure was measured in the experiments, which showed that the FeCuNbSiB/piezofiber structure oscillated during the natural period $T = \sim 37.78 \, \mu s$ in the time domain. Specifically, the self-vibration phenomena vanished after the excited field vanished when the width was equal to the multiples of the natural period $T$. The transient ME response was strongly dependent on the external bias magnetic field, which was the same as with the ME response driven by
the sine magnetic field. The sensing characteristics to the pulsed magnetic field were also measured. The results show that the output sensitivities reached 17 mV/Oe Δt = 100 µs for detecting the amplitude of the pulse magnetic field and reached 5.4 mV/µs when Δt = 0–300 µs for detecting the pulsed width. The results obtained may be useful for the development of ME sensors for pulsed magnetic fields, such as the lighting current monitoring field.

**Author Contributions:** C.L. designed the project; H.Z. developed ME composite and carried out experimental works, A.Y. designed the experimental system. Z.O. carried out simulation work. F.Y. performed the experimental works. H.G. performed experiments and analyzed results; C.L. and H.Z. wrote the manuscript.

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