INTRODUCTION

As polymer networks infiltrated with water, hydrogels have been widely used as scaffolds for tissue engineering (1), vehicles for drug delivery (2), and model platforms for biological studies (3). More recently, hydrogels have been explored for applications in devices and machines, including wearable electronics (4, 5), stretchable optical fibers (6), adhesives (7–9), soft robotics (10), and hydrogel-based soft machines (11, 12). The use of hydrogels in devices and machines requires them to maintain robustness under cyclic mechanical loads. Following the pioneering work by Gong et al. (13), hydrogels have been made tough to resist crack propagation under a single cycle of mechanical load (14–17). The toughening of hydrogels is achieved by integrating mechanisms for dissipating mechanical energy such as fracture of short polymer chains and reversible cross-links into stretchy polymer networks (18, 19). However, existing tough hydrogels still suffer from fatigue fracture under multiple cycles of mechanical loads (11, 20, 21), because the resistance to fatigue crack propagation after prolonged cycles of loads is the energy required to fracture a single layer of polymer chains (i.e., the intrinsic fracture energy of the hydrogel), which is unaffected by the additional dissipation mechanisms introduced in tough hydrogels (20, 21). The reported fatigue thresholds of various tough hydrogels are 8.4 J/m² for polyacrylamide (PAAm)–polyvinyl alcohol (PVA) (21) and 53.2 J/m² for PAAm-alginate (20), the same order as their intrinsic fracture energies. The highest fatigue threshold for hydrogels reported so far is 418 J/m² for a double-network hydrogel, poly(2-acrylamido-2-methyl-1-propanesulfonic acid) (PAMPS)–PAAm, which possibly can be attributed to the PAAm network with very long polymer chains, and thus, high intrinsic fracture energy (13, 22). A general strategy toward the design of anti-fatigue-fracture hydrogels has remained a critical need and a central challenge for long-term applications of hydrogels in devices and machines.

In contrast to synthetic hydrogels, biological tissues such as cartilages, tendons, muscles, and heart valves show extraordinary anti-fatigue properties. For example, the knee joint of an average person needs to sustain peak stresses of 4 to 9 MPa for 1 million cycles per year, and its fracture energy after prolonged cycles of loads is above 1000 J/m² (23, 24). The anti-fatigue property of biological tissues possibly arises from the inherently highly ordered and partially crystalline structures of collagen fibers in the tissues (25). Inspired by anti-fatigue biological tissues, here, we hypothesize that the increase of crystallinity in synthetic hydrogels can substantially enhance their anti-fatigue-fracture properties. The fatigue threshold of polyvinyl alcohol (PVA) with a crystallinity of 18.9% weight % in the swollen state can exceed 1000 J/m².

RESULTS

To validate the hypothesis that the increase of crystallinity in hydrogels can greatly enhance their fatigue thresholds, we use PVA hydrogels as a model material system with tunable crystallinity. We first freeze a solution of uncross-linked PVA at −20°C for 8 hours and thaw it at 25°C for 3 hours to form a hydrogel cross-linked by crystalline domains (27). The freeze-thawed PVA hydrogel is further dried in an incubator at 37°C and then annealed at 100°C for various times ranging from 0 to 90 min (28). (The hydrogel dry-annealed for 0 min means no annealing process.) The crystallinity of the hydrogel can be tuned by drying and...
substantially more crystalline domains nucleate during the air-drying process (see Fig. 2A). The increase of the crystallinity implies that, when the freeze-thawed PVA (without excess chemical cross-links) is dried in air, the crystallinity in the dry state increases to 2.1 wt %, respectively. How-ever, when the freeze-thawed PVA (without excess chemical cross-links) is dried in air, the crystallinity in the dry state increases to 37.7 wt % (see Materials and Methods for details). The excess chemical cross-links minimize the formation of further crystalline domains during the air-drying process (29, 30). As shown in Fig. 2A, both the chemically cross-linked and freeze-thawed PVA hydrogels show negligible endothermic peaks, with measured crystallinities in the dry state of 0.2 and 2.1 wt %, respectively. However, when the freeze-thawed PVA (without excess chemical cross-links) is dried in air, the crystallinity in the dry state increases to 37.7 wt % (see Fig. 2A). The increase of the crystallinity implies that substantially more crystalline domains nucleate during the air-drying process (28). The crystallinity further increases gradually by increasing the annealing time. When the sample is annealed for 90 min at 100°C, the crystallinity in the dry state reaches 47.3 wt %. In addition to the crystallinities in the dry state measured from DSC, we further measure the water contents in the fully swollen samples shown in Fig. 2B and calculate the corresponding crystallinities in the swollen state shown in Fig. 2C. It is well known that the as-prepared dry PVA samples may contain residual water bonded with polymer chains (30). The amount of residual water can be calculated from the endothermic transition ranging from 60° to 180°C on the DSC curves (see Fig. 2A and fig. S1) (27). The above reported crystallinities in the dry state (without residual water), water contents, and crystallinities in the swollen state have been corrected to account for the weights of residual water in the as-prepared dry samples (see Materials and Methods, table S1, and fig. S1 for details).

To quantify the evolution of crystalline morphology, we measure the average distance between adjacent crystalline domains $L$ through small-angle x-ray scattering (SAXS) and the average size of crystalline domains $D$ through wide-angle x-ray scattering (WAXS) (see details in Materials and Methods). We first perform SAXS measurements on the samples in the swollen state after subtracting the water background, measuring the scattering intensity $I(q)$ versus the scattering vector $q$. To identify the location of the peak intensity, we correct the intensity by multiplying the scattering intensity with the scattering vector $q$ (31). As shown in Fig. 2D, there is no peak in the plot of the corrected intensity $I_q$ versus the scattering vector $q$ for the freeze-thawed hydrogel, which implies negligible interference between adjacent crystalline domains. For the hydrogel dry-annealed for 0 min, there is a slight shoulder shown in the corrected intensity curve, which indicates stronger interference between adjacent domains. The average distance between adjacent crystalline domains $L$ can be estimated from the critical vector corresponding to the peak intensity $q_{\text{max}}$ following the Bragg expression $L = 2\pi/q_{\text{max}}$ (31). The average distance for the hydrogel dry-annealed for 0 min is estimated to be 21 nm in the swollen state. As the annealing time increases to 90 min, the average distance $L$ decreases to 13 nm in the swollen state (Fig. 2G). As a control case, we also measure SAXS profiles of the hydrogel dry-annealed for 90 min in the dry state. As shown in Fig. 2F, the average distance between adjacent crystalline domains in the dry state is around 9 nm, smaller than the distance in the...
same hydrogel in the swollen state (i.e., 13 nm). This is because swelling of the interstitial amorphous polymer chains increases the distance between adjacent crystalline domains.

We further perform WAXS measurements on the hydrogels in their dry state using Ni-filtered CuK\textsubscript{α} radiation with x-ray wavelength \(l = 1.54\) Å. As shown in Fig. 2E, all dry-annealed PVA hydrogels show a strong diffraction peak at \(2θ = 19.7°\), which corresponds to the typical reflection plane of (10\(\overline{1}\)) in semicrystalline PVA. Small peaks at \(2θ = 11.5°\) and 23.1° are also observed in the hydrogel dry-annealed for 90 min, suggesting a high crystallinity in the hydrogel, which is consistent with the DSC measurement. By identifying the half width of the maximum diffraction peak \(b\), the average size of crystalline domains \(D\) can be approximately calculated using Scherrer’s equation \(D = k\lambda/(β \cos θ)\), where \(k\) is a dimensionless shape factor varying with the actual shape of the crystalline domain, \(λ\) is the wavelength of x-ray diffraction, and \(θ\) is the Bragg angle. Here, \(β\) is identified after subtracting the instrumental line broadening, and the dimensionless shape factor \(k\) is set as 1, approximating the spherical shape of the crystalline domains.

As shown in Fig. 2G, by increasing the annealing time from 0 to 90 min, the average size of the crystalline domains increases from 3.8 to 6.5 nm. This trend is consistent with the decrease of the average distance between adjacent crystalline domains with annealing time, because the growth of the crystalline domains consumes the interstitial amorphous polymer chains.

To further validate the tuning of crystalline domains in the PVA hydrogel with annealing time, we use tapping-mode atomic force microscopy (AFM) to obtain phase images of the hydrogels dry-annealed for 0 and 90 min, respectively. The bright areas in Fig. 2H correspond to the regions with relatively high modulus (mainly crystalline domains), whereas the dark areas represent the regions with relatively low modulus (mainly amorphous domains). As shown in Fig. 2H, the morphology of isolated crystalline domains is observed in the hydrogel dry-annealed for 0 min, while the hydrogel dry-annealed for 90 min shows larger aggregated crystalline domains.

**Characterization of fatigue fracture properties of hydrogels**

To measure the fatigue threshold of PVA hydrogels, we adopt the single-notch method, which is widely used in fatigue tests of rubbers. Notably, all fatigue tests in this study are performed on fully swollen hydrogels immersed in a water bath to prevent the dehydration-induced crack propagation (fig. S2). We use dogbone-shaped samples and perform cyclic tensile tests on both notched and unnotched samples,
which are otherwise the same (Fig. 1, B and C). The nominal stress versus stretch curves (i.e., \( S \) versus \( \lambda \)) of unnotched samples are obtained over \( N \) cycles of applied stretch \( \lambda^A \). The strain energy density \( W \) of the unnotched sample under the \( N \)th cycle of applied stretch \( \lambda^A \) can be calculated as

\[
W(\lambda^A, N) = \int_1^{\lambda^A} Sd\lambda
\]

(1)

where \( S \) and \( \lambda \) are the measured nominal stress and stretch, respectively. Thereafter, the same cyclic stretch \( \lambda^A \) is applied on the notched sample, measuring the evolution of the cut length in undeformed state \( c \) with the cycle number \( N \). The applied energy release rate \( G \) in the notched sample under the \( N \)th cycle of applied stretch \( \lambda^A \) can be calculated as (35, 36)

\[
G(\lambda^A, N) = 2k(\lambda^A) \cdot c(N) \cdot W(\lambda^A, N)
\]

(2)

where \( k \) is a slowly varying function of the applied stretch expressed as \( k = 3/\sqrt{\lambda^A} \), \( c \) is the current crack length at undeformed configuration, and \( W \) is the strain energy density measured in the unnotched sample (Eq. 1). By systematically varying the applied stretch \( \lambda^A \), we can obtain a plot of crack extension per cycle versus the applied energy release rate (i.e., \( dc/dN \) versus \( G \)). By linearly extrapolating the curve of \( dc/dN \) versus \( G \) to the intercept with the abscissa, we can approximately obtain the critical energy release rate \( G_c \), below which the fatigue crack will not propagate under infinite cycles of loads. By definition, the fatigue threshold \( G_\text{f} \) is equal to the critical energy release rate \( G_c \). To validate that this extrapolated value \( G_c \) is the fatigue threshold \( G_\text{f} \), we further apply \( G_c \) to the notched sample over 30,000 cycles (to approximate infinite cycles of loads) and observe no crack extension (Fig. 3H). In addition, we also measure the fatigue threshold using the pure-shear test (20) to validate the results from the single-notch test (Fig. 6).

For cyclic tensile tests on unnotched samples (Fig. 3, A and B, and fig. S3), both chemically cross-linked PVA and freeze-thawed PVA show negligible Mullins effect, and their \( S \) versus \( \lambda \) curves reach steady states after only a few cycles (i.e., 10 for chemically cross-linked PVA at an applied stretch of 1.6 and 200 for freeze-thawed PVA at an applied stretch of 2.2). On the other hand, the dry-annealed PVA hydrogels exhibit a more significant Mullins effect due to mechanical dissipation caused by melting and reorientation of crystalline domains (31). As the stretch further increases, the crystalline domains may transform into aligned fibrils along the loading direction (31). The energy required to damage the crystalline domains and fibrils is much higher than that to fracture a single layer of the same polymer in amorphous state. The hydrogel dry-annealed for 90 min reaches a steady state after 1000 cycles of applied stretches of \( \lambda^A = 2 \) (Fig. 3C). Despite the Mullins effect, the steady-state maximum nominal stress of the hydrogel dry-annealed for 90 min is much higher than that of both the chemically cross-linked PVA and the freeze-thawed PVA at the same applied stretch of \( \lambda^A = 2 \) (e.g., 2.3 MPa for hydrogel dry-annealed for 90 min, 15 kPa for chemically cross-linked PVA, and 3 kPa for freeze-thawed PVA in fig. S4).

For cyclic tensile tests on notched samples, a pre-crack is cut using a razor blade with a tip radius of around 200 \( \mu \)m and an initial crack length of around 1 mm, smaller than one-fifth of the width of the sample (36). A digital microscope (AM48152ZT, Dino-Lite; resolution, 20 \( \mu \)m/pixel) is used to record the cut length under cyclic loads (Materials and Methods and fig. S5). We first apply cyclic loads with a small applied stretch (i.e., \( \lambda^A = 1.3 \)) on a notched sample. If the crack remains quasi-stationary with crack extension per cycle (\( dc/dN \)) smaller than 20 nm per cycle (i.e., no detectable crack extension in 1000 cycles), the applied cyclic stretch is increased by the increment of \( \Delta \lambda^A = 0.1 \) for other notched samples until crack propagation greater than 20 nm per cycle is captured. As shown in Fig. 3 (D and E), the fatigue thresholds of chemically cross-linked PVA and freeze-thawed PVA are measured to be 10 and 23 J/m\(^2\), respectively. The fatigue threshold for the hydrogel dry-annealed for 0 min increases to 110 J/m\(^2\). As the annealing time increases, the fatigue threshold further increases. In particular, for the hydrogel dry-annealed for 90 min, the fatigue threshold can achieve 1000 J/m\(^2\) (Fig. 3F and fig. S5). The measured fatigue threshold of the hydrogel dry-annealed for 90 min from the pure-shear test is 918 J/m\(^2\), consistent with the single-notch test (fig. S6). The dependence of fatigue threshold on the crystallinity is summarized in Fig. 3G. The fatigue threshold increases with the crystallinity and demonstrates a sharp jump when the crystallinity in the swollen state reaches approximately 15 wt %.

In addition to fatigue tests, we also measure the nominal stress versus stretch curves of all hydrogels to obtain their Young’s moduli and tensile strengths (fig. S7). As shown in Fig. 4 (A and B), both the Young’s modulus and tensile strength increase with the hydrogels’ crystallinity and show marked enhancements when the crystallinity in the swollen state reaches approximately 15 wt % (37). This sharp jump in Young’s modulus and tensile strength is consistent with the marked increase of fatigue threshold of the hydrogel at approximately 15 wt % crystallinity in the swollen state (Fig. 3G).

**Patterning of highly crystalline regions in hydrogels**

Whereas annealing the whole PVA hydrogel can greatly enhance its fatigue threshold, the annealing treatment also increases the modulus and decreases the water content of the hydrogel (Fig. 4C and fig. S8). However, for many applications, it is desirable to maintain the relatively low modulus and high water content of the hydrogel. Here, we propose a strategy to introduce programmed highly crystalline regions in the hydrogels. We use the computer-aided design of electrical circuits to induce localized heat treatment for annealing selected regions of the hydrogels (see details in Materials and Methods and fig. S9). The chemically cross-linked PVA hydrogel is used as the pristine sample with low fatigue threshold of 15 J/m\(^2\), low Young’s modulus of 114 kPa, and high water content of 88 wt %. Three examples of programmed patterning patterns on the pristine PVA hydrogels are demonstrated with enhanced fatigue thresholds.

We first locally introduce a highly crystalline ring-shaped region around a crack tip (Fig. 5A), leading to a fatigue threshold of more than 236 J/m\(^2\) (Fig. 5B). Despite the small area of the highly crystalline region, this local conditioning can greatly reinforce the crack tip, delaying crack propagation. Meanwhile, the measured overall Young’s modulus and water content of the sample are maintained at 114 kPa and 87 wt % (Fig. 5, E and F), respectively, which are almost unaffected by the local annealing around the crack tip. As a second example, we pattern mesh-like highly crystalline regions on the pristine hydrogel (Fig. 5C). Compared with the pristine sample, the fatigue threshold of the mesh-reinforced sample increases to 290 J/m\(^2\) (Fig. 5D), and its Young’s modulus remains relatively low (627 kPa) and its water content stays relatively high (83 wt %) (Fig. 5, E and F).

The strategy of patterning highly crystalline regions can be applied to various structures of hydrogels for improving anti-fatigue performances as well. For example, the kirigami structure is commonly adopted to...
enhance the stretchability (38) and effective adhesion (39) of films. As a third example, we demonstrate a kirigami hydrogel sheet with improved anti-fatigue performance by patterning highly crystalline regions around the cut tips. As shown in Fig. 5G, we first generate a kirigami pattern on a pristine sample by introducing parallel periodic cuts with equal length and equal distance between adjacent cuts. The ultimate stretch and effective nominal stress (i.e., force divided by cross-sectional area) of the pristine kirigami hydrogel sheet under a single cycle of load are measured to be 1.7 and 0.9 kPa, respectively. Thereafter, the pristine kirigami sheet is reinforced around the cut tips with local annealing. The reinforced kirigami sheet can sustain cyclic tensile loads at an applied stretch of \( \lambda^A = 2.1 \) without detectable fatigue fracture even when the cycle number approaches 3000. Moreover, the effective nominal stress after prolonged cycles at the applied stretch of \( \lambda^A = 2.1 \) can maintain the plateau of 13 kPa, which is 14.4 times the strength of the pristine kirigami sheet (Fig. 5, H and I).

Figure 5 (J and K) compares fatigue thresholds, water contents, and Young's moduli of reported hydrogels in literature (20–22, 40, 41). We
**Fig. 5. Patterning highly crystalline regions in PVA hydrogels.** (A) Illustration of introducing a highly crystalline region around crack tip. Inset: Raman spectroscopy with bright color representing low water content and dark color representing high water content (see details in Materials and Methods). (B) Comparison of crack extension per cycle dc/dN versus applied energy release rate G between the pristine sample and the tip-reinforced sample. The fatigue thresholds of the pristine sample and the tip-reinforced sample are 15 and 236 J/m², respectively. (C) Illustration of introducing mesh-like highly crystalline regions. Inset: Digital image correlation (DIC) method shows large deformation in low-crystallinity regions and small deformation in high-crystallinity regions. (D) Crack extension per cycle dc/dN versus applied energy release rate G of the mesh-reinforced sample. The fatigue threshold of the mesh-reinforced sample is 290 J/m². (E) Water contents of the pristine sample, the tip-reinforced sample, the mesh-reinforced sample, and the fully annealed sample. (F) Young’s moduli among reported synthetic hydrogels, PVA hydrogels with patterned highly crystalline regions, and biological tissues. (G) Illustration of introducing highly crystalline regions around cut tips in a pristine kirigami sheet. (H) Effective nominal stress versus stretch curves of the reinforced kirigami sheet under cyclic loads. Effective nominal stress versus stretch curve of the pristine kirigami sheet under a single cycle of load. (I) Images of the reinforced kirigami sheet under 1000th cycle and under 3000th cycle. (J) Comparison of fatigue thresholds and water contents among reported synthetic hydrogels (20–22, 40, 41), PVA hydrogels with patterned highly crystalline regions, and biological tissues (23). (K) Comparison of fatigue thresholds and Young’s moduli among reported synthetic hydrogels, PVA hydrogels with patterned highly crystalline regions, and biological tissues. IPN in (J) and (K) represents interpenetrating polymer network. Data in (E) and (F) are means ± SD, n = 3. Scale bars, 800 μm (A), 1 mm (C), and 40 mm (I).
show that by patterning highly crystalline regions, both tip- and mesh-reinforced PVA hydrogels outperform existing synthetic hydrogels in terms of fatigue thresholds, and they can still maintain relatively high water contents and relatively low Young’s moduli.

**DISCUSSION**
We have proposed that the design principle for anti-fatigue-fracture hydrogels is to make the fatigue crack encounter fracture objects requiring energies per unit area much higher than that for fracturing a single layer of amorphous polymer chains. We have demonstrated that the fatigue thresholds of hydrogels can be greatly enhanced by designing crystalline domains in the hydrogels. We use PVA hydrogels as a model material to validate this new mechanism for designing anti-fatigue-fracture hydrogels. The fatigue threshold of a PVA hydrogel with a crystallinity of 18.9 wt % in the swollen state can achieve more than 1000 J/m². We further develop a strategy to make anti-fatigue-fracture PVA hydrogels but still maintain their high water contents and low moduli by patterning highly crystalline regions in the hydrogels. The reported mechanism and strategy for designing anti-fatigue-fracture hydrogels can be extended to hydrogel composites with fillers such as nanoclay, nanocellulose, and nanofibers.

The capability to enhance the anti-fatigue-fracture performance of synthetic hydrogels makes a number of future research directions and applications possible. For example, anti-fatigue-fracture hydrogels can be used for hydrogel-based gastric-retentive devices and implantable tissue replacements of meniscus, intervertebral disc, and cartilage, which require long-term mechanical robustness when interacting with the human body.

**MATERIALS AND METHODS**

**Synthesis of PVA hydrogels**
All types of our PVA hydrogels were synthesized from 10 wt % PVA (molecular weight, 146,000 to 186,000; 99+% hydrolyzed; Sigma-Aldrich, 363065) solution. The solution was heated in a water bath at 100°C with stirring for 5 hours. To synthesize chemically cross-linked PVA hydrogels, we added 10 μl of glutaraldehyde [25 volume percent (volume %); Sigma-Aldrich, G6257] as a cross-linker to 1 ml of 10 wt % PVA solution and added 10 μl of hydrochloric acid (36.5 to 38 wt %; J.T. Baker, 9535-02) as an accelerator into the other 1 ml of 10 wt % PVA solution. We then mixed and defoamed each of them by using a centrifugal mixer (AR-100, Thinky). The final mixtures, obtained by mixing and defoaming the two solutions together, were then cast into a mold and allowed to cure for 2 hours. The chemically cross-linked PVA hydrogels were immersed in deionized (DI) water for 2 days to remove unreacted chemicals. To fabricate freeze-thawed PVA hydrogels, 10 wt % PVA solutions after mixing and defoaming were poured into a mold, frozen at −20°C for 8 hours, and thawed at 25°C for 3 hours. The freeze-thawed hydrogels were further dried in an incubator (New Brunswick Scientific, C25) at 37°C for 2 hours and annealed at 100°C for a controlled time (i.e., 0, 1, 3, 5, 10, or 90 min). All as-prepared PVA hydrogels were immersed in water to achieve their equilibrium-swollen state.

**Measurement of residual water and crystallinity in dry samples**
We measured the crystallinities of the resultant PVA hydrogels in their dry state by DSC (DSC/cell: RCS1-3277 Cooling System: DSC1-0107).

For as-prepared chemically cross-linked PVA and freeze-thawed PVA, we used excess chemical cross-links to fix the amorphous polymer chains before air-drying, minimizing the formation of further crystalline domains during the air-drying process. We first soaked the samples (thickness of 1 mm) in the aqueous solution consisting of 10 ml of glutaraldehyde (25 volume %; Sigma-Aldrich, G6257), 500 μl of hydrochloric acid (36.5 to 38 wt %; J.T. Baker, 9535-02), and 50 ml of DI water for 2 hours. Thereafter, we soaked the samples in a DI water bath for 2 hours to remove the residual hydrochloric acid. The samples were further dried in an incubator (New Brunswick Scientific, C25) at 37°C for 2 hours. Thereafter, we measured the mass of residual water \( m_{\text{residual}} \) and the total mass of crystalline domains \( m_{\text{crystalline}} \) and the total mass of the dry samples (with residual water) \( m \) using DSC. In a typical DSC measurement, we first weighed the total mass of the dry sample (with residual water) \( m \). The sample was thereafter placed in a Tzero Pan and heated up from 50°C to 250°C at a rate of 20°C/minute under a nitrogen atmosphere with a flow rate of 30 ml/min. The curve of heat flow shows a broad peak from 60°C to 180°C, indicating that the sample contains a small amount of residual water. The integration of the endothermic transition ranging from 60°C to 180°C gives the enthalpy for evaporation of the residual water per unit mass of the dry sample (with residual water) \( H_{\text{water}} \). Therefore, the mass of the residual water \( m_{\text{residual}} \) can be calculated as

\[
    m_{\text{residual}} = m - \frac{H_{\text{water}}}{H_{\text{crystalline}}} \tag{3}
\]

where \( H_{\text{water}} = 2260 \text{J/g} \) is the latent heat of water evaporation. The curve of heat flow shows another narrow peak ranging from 200°C to 250°C, corresponding to melting of the crystalline domains. The integration of the endothermic transition ranging from 200°C to 250°C gives the enthalpy for melting the crystalline domains per unit mass of the dry sample (with residual water) \( H_{\text{crystalline}} \). Therefore, the mass of the crystalline domains \( m_{\text{crystalline}} \) can be calculated as

\[
    m_{\text{crystalline}} = m - \frac{H_{\text{crystalline}}}{H_{\text{crystalline}}} \tag{4}
\]

where \( H_{\text{crystalline}} = 138.6 \text{J/g} \) is the enthalpy of fusion of 100 wt % crystalline PVA measured at the equilibrium melting point \( T_m^{\text{dry}} \). Therefore, the crystallinity in the ideally dry sample \( X_{\text{dry}} \) (without residual water) can be calculated as

\[
    X_{\text{dry}} = \frac{m_{\text{crystalline}}}{m - m_{\text{residual}}} \tag{5}
\]

**Measurement of water content and crystallinity in swollen samples**
The swollen hydrogels weighing \( m_{\text{swollen}} \) were placed in an incubator (New Brunswick Scientific, C25) at 37°C for 2 hours and weighed \( m \) after air-drying. The mass of the residual water in the as-prepared dry samples \( m_{\text{residual}} \) was measured in the previous section. Therefore, the water content in the swollen state can be calculated as \( (m_{\text{swollen}} - m + m_{\text{residual}})/m_{\text{swollen}} \) and the polymer content in the swollen state can be calculated as \( (m - m_{\text{residual}})/m_{\text{swollen}} \). In addition, the crystallinity in the swollen sample \( X_{\text{swollen}} \) can be calculated as

\[
    X_{\text{swollen}} = \frac{m_{\text{crystalline}}}{m_{\text{swollen}}} = X_{\text{dry}} \cdot \frac{m - m_{\text{residual}}}{m_{\text{swollen}}} \tag{6}
\]
where \( m_{\text{crystalline}} \) and \( X_{\text{dry}} \) were measured as described in the previous section.

### AFM phase imaging

AFM phase images were acquired with an atomic force microscope (MFP-3D, Asylum Research) in tapping mode. Dry freestanding PVA films were directly attached onto the sample stage with a double-sided carbon tape. The probe lightly taps on the sample surface with a recorded phase shift angle of the probe motion relative to a driving oscillator. The bright regions with high phase angle correspond to regions with relatively high modulus, and the dark regions with low phase angle correspond to regions with relatively low modulus.

### X-ray scattering

The x-ray scattering measurement was performed with a Pilatus3R 300K detector [Bruker Nanostar SAXS in X-ray Diffraction Shared Experimental Facility at Massachusetts Institute of Technology (MIT)]. We used a small-angle 2-mm beamstop with a sample-detector distance of 1059.1 mm for SAXS measurements and a wide-angle 2-mm beamstop with a sample-detector distance of 109.1 mm for WAXS measurements. The exposure time was set as 300 s. Raw SAXS and WAXS patterns were processed with corrections by MATLAB-based GIXSGUI software before analysis.

### Raman spectroscopy

Each sample was hydrated for more than 2 hours and then pressed between a glass slide and a coverslip to ensure a flat surface. The coverslip was then sealed at the edges with nail polish to prevent the hydrogel from drying. A confocal Raman microscope (alpha300 RA, WITec, Germany) with a 20× objective (Zeiss, Germany) was used. An Nd:YAG laser (532 nm) was used as the excitation source with a maximum power of 75 mW. Data were collected using a charge-coupled device detector (DU401A, Andor, UK) behind a grating spectrometer (600 g/mm; UHTS 300, WITec, Germany). A 20-μm-resolution Raman map of 4 × 3-mm scan area was acquired with an accumulation time of 1 s per point. Each point was prebleached for 400 ms to decrease the effect of fluorescence. Cosmic ray removal and background subtraction were performed to clean the spectra. The intensity of O–H bond within the PVA and water was calculated by integrating the spectra in the range of 2800 to 3000 cm\(^{-1}\) and 3075 to 3625 cm\(^{-1}\), respectively. The ratio of PVA and water was then calculated and plotted as a heatmap shown in Fig. 5A.

### Measurement of the fatigue threshold

All the mechanical tests were performed in a water bath at 25°C with a U-stretch testing device (CellScale, Canada). For mechanically weak samples (e.g., the chemically cross-linked and freeze-thawed hydrogels), a load cell with a maximum force of 4.4 N was used; for mechanically strong samples (e.g., hydrogel dry-annealed for 90 min), a load cell with a maximum force of 44 N was used. The dogbone-shaped sample had the dimensions at as-prepared state with a width of 5 mm, a thickness of 0.8 mm, and a gauge length of 10 mm. The nominal stress \( S \) was measured from the recorded force \( F \) divided by width \( W \) and thickness \( t \) in the swollen state. To measure the applied stretch \( \lambda^{\text{A}} \) in gauge length, we adopted the digital image correlation method and calibrated the correlation between \( \lambda^{\text{A}} \) and loading distance \( d \) at gripping points (fig. S2). A digital microscope (AM4815ZT, Dino-Lite; resolution, 20 μm/pixel) was used to record the crack extension. Because dry-annealed PVA was transparent when immersing in water, we spread a small amount of graphite powder on the surface of the sample for visualization.

### Preparation of samples with programmable crystalline domains

A copper tape was first loaded in a paper-cutting machine (Silhouette CAMEO 3). The electrical circuits were first designed with computer-aided design and loaded to the machine via Silhouette Studio software. After completion of the cutting, the copper tape was placed on a flat acrylic plate and unwanted parts were carefully peeled off. Two additional conductive wires were soldered on the copper tape as two electrodes. An electrical generator (Instek PSR-2400L2, T’Equipe) was used to apply controlled current around 2 to 5 mA on the designed circuits. To ensure that the designed temperature was achieved, a thermal imager (Seek Thermal XR imager) was first used to measure the local temperature from the circuit (see fig. S9). By adjusting the current to control the design temperature, the air-dried PVA hydrogel was placed on the targeted circuits for local heating. A gripper was used to ensure full contact between the PVA films and copper circuits.

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