Structure and insertion properties of disordered and ordered LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinels prepared by wet chemistry

Abstract We present the synthesis, characterization, and electrode behavior of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinels prepared by the wet-chemical method via citrate precursors. The phase evolution was studied as a function of nickel substitution and upon intercalation and deintercalation of Li ions. Characterization methods include X-ray diffraction, SEM, Raman, Fourier transform infrared, superconducting quantum interference device, and electron spin resonance. The crystal chemistry of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ appears to be strongly dependent on the growth conditions. Both normal-like cubic spinel \([Fd\bar{3}m\text{ space group (SG)}]\) and ordered spinel \((P4_132\text{ SG})\) structures have been formed using different synthesis routes. Raman scattering and infrared features indicate that the vibrational mode frequencies and relative intensities of the bands are sensitive to the covalency of the \((\text{Ni, Mn})-\text{O}\) bonds. Scanning electron microscopy (SEM) micrographs show that the particle size of the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ powders ranges in the submicronic domain with a narrow grain-size distribution. The substitution of the 3d$^8$ metal for Mn in LiNi$_{0.5}$Mn$_{1.5}$O$_4$ oxides is beneficial for its charge–discharge cycling performance. For a cut-off voltage of 3.5–4.9 V, the electrochemical capacity of the Li/LiNi$_{0.5}$Mn$_{1.5}$O$_4$ cell is ca. 133 mAh/g during the first discharge. Differences and similarities between LiMn$_2$O$_4$ and LiNi$_{0.5}$Mn$_{1.5}$O$_4$ oxides are discussed.

Keywords Spinel · Lithium batteries · Nickel doping · Wet synthesis · Local structure

Introduction

The cyclability of Li$^+$ ion transfer in Li-ion rechargeable batteries depends mainly on the dimensional stability of the host material during insertion and deinsertion of Li$^+$ [1]. Recently, it was found that redox reactions of nonstoichiometric or doped Li$_{x}$Mn$_{2-y}$O$_4$ (\(M=\text{Li, Co, Cr, Ni, Al, etc.}\)) spinels are much better than that of the pure LiMn$_2$O$_4$ ceramics, which exhibit a slight capacity fading [2]. These materials have been investigated to improve the cycling performance of LiMn$_2$O$_4$, showing a tetragonal distortion induced by the excess of Mn$^{3+}$ Jahn–Teller ions in deeply discharged electrode. However, all the reported doping methods have led to a decreased specific charge compared to the undoped LiMn$_2$O$_4$ materials so far. The search of high-voltage material electrodes has been focused on two categories: the inverse spinels, e.g., LiMVO$_4$, and the normal spinels, Li$_{x}$M$_{2-x}$Mn$_2$O$_4$.

Recent investigations have shown that, among the Ni-substituted LiMn$_2$O$_4$ spinels, the composition LiNi$_{0.5}$Mn$_{1.5}$O$_4$ possesses specific electrochemical characteristics such as a high capacity of 130–140 mAh/g associated with a high-voltage plateau in the 5-V range [3–25]. Gao et al. [3] investigated the origin of the voltage profile for LiNi$_{x}$Mn$_{2-x}$O$_4$. For the same system, Zhong et al. [4] showed the effects of the synthesis route (sol-gel vs solid state) on some structural and electrochemical properties. The almost flat voltage profile was confirmed by Ohzuku et al. [13]. All these works only reported on structural and electrochemical information, and few of them
concerned the physical properties of the LiNi_{0.5}Mn_{1.5}O_{4} material. Blasse [26] showed in 1966 that LiNi_{0.5}Mn_{1.5}O_{4} is a ferrimagnetic spinel material with Ni^{2+} and Mn^{4+} on octahedral sites and Li on tetrahedral sites. In their studies of LiNi_{0.5}Mn_{1.5}O_{4} by neutron diffraction and IR spectroscopy, Gryffroy and Vandenberghe [27] showed in terms of the 1:3 cation-ordering at B sites of the A[B_{2}]O_{4} spinel structure, i.e., 16d sites in Fd3m, with a cubic symmetry P4_{3}3_{2} space group (SG). A schematic representation of the structure of the normal spinel and the 1:3 ordered spinel is shown in Fig. 1. Recently, the formation of the ordered spinel has been confirmed by combining X-ray diffraction (XRD) and Fourier transform infrared (FTIR) measurements [6, 15].

Several methods have been used to synthesize LiNi_{0.5}Mn_{1.5}O_{4}. These methods are solid-state reaction [7], sol-gel method [8, 9], molten salt technique [10], emulsion drying method [11], composite carbonate process [12], coprecipitation method [13], and melt-impregnation synthesis [14]. Thin-films of LiNi_{0.5}Mn_{1.5}O_{4} have been prepared by electrostatic spray deposition [15]. Results have shown that LiNi_{0.5}Mn_{1.5}O_{4} materials were obtained with different properties. In fact, the key parameter in electrochemical capacity fading could be rather the oxide morphology (grain size and/or specific surface area), as shown by the remarkable improvement reported on sol-gel LiNi_{0.5}Mn_{1.5}O_{4} [25].

The aim of this communication is the growth and characterization of the LiNi_{0.5}Mn_{1.5}O_{4} materials prepared by wet chemistry. We use two different techniques, i.e., the sol-gel method assisted by citric acid and the combustion technique assisted by glycine. The LiNi_{0.5}Mn_{1.5}O_{4} system crystallizes, in peculiar synthesis conditions, as either an ordered- or a normal-spinel phase. The synthesis process has been studied by thermogravimetry (TG) and differential thermal analysis (DTA). The long-range structure and local atomic environment are investigated by XRD, Raman scattering (RS), and FTIR spectroscopy. Electronic properties are investigated by electron spin resonance (ESR) and superconducting quantum interference device (SQUID) magnetometry, which also probe the local structure. Finally, we also report the difference in electrochemical properties of the two structural types of LiNi_{0.5}Mn_{1.5}O_{4} materials.

**Experimental description**

**Synthesis**

To compare the effect of the synthesis method, two LiNi_{0.5}Mn_{1.5}O_{4} samples were prepared by sol-gel synthesis and combustion process, both using carboxylic acid as chelating agent. The choice of these wet-chemical techniques has been guided by the remarks of Zhong et al. [4]. The sol-gel and combustion synthesis occurred from metal acetates via inorganic polymerization reactions in solution according to methods reported in our previous works [28, 29]. Compared with the conventional solid-state reaction, it is believed that these techniques offer many advantages such as lower temperature processing and better control of material morphology.

The procedure for the sol-gel synthesis has been detailed elsewhere [28]. Reagent grade Li(CH_{3}COO)·2H_{2}O, Ni(CH_{3}COO)_{2}·4H_{2}O, and Mn(CH_{3}COO)_{2}·4H_{2}O were selected as raw materials. The powder samples of LiNi_{0.5}Mn_{1.5}O_{4} were prepared using stoichiometric amounts of Li(CH_{3}COO)·2H_{2}O, Ni(CH_{3}COO)_{2}·4H_{2}O, and Mn(CH_{3}COO)_{2}·4H_{2}O acetate dissolved in distilled water and added to a continuously stirred aqueous solution, and then evaporated at 80 °C until a transparent gel was obtained. The resulting slurries were heated at 400 °C for a few hours to get the product of the composition LiNi_{0.5}Mn_{1.5}O_{4}. Accordingly, the powder mass was slightly ground and then fired at 680 °C in air for 5 h to improve the crystallinity of the LiNi_{0.5}Mn_{1.5}O_{4} final product.

The combustion method (glycine-nitrate pyrolysis) has been reported by Chitra et al. and Julien et al. [29, 30]. In this technique, stoichiometric amounts of Li-, Ni-, and Mn-nitrate salts were dissolved in triple-distilled water with a small proportion of ethanol and aqueous solution of glycine (amino-acetic acid, NH_{2}CH_{2}COOH). In this case, care was exercised in adjusting the concentration of the complexing agent; the pH of the solution ranged between 5 and 6.

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**Fig. 1** Schematic representation of the structure of AB_{2}O_{4} spinel lattices. The smallest (primitive) cubic unit cell of normal spinel (Fd3m) and the unit cell of the 1:3 ordered spinel (P4_{3}3_{2}). The structure is composed of alternating octants of AO_{4} tetrahedra and B_{2}O_{4} cubes to build the face-centered cubic unit cell.
Apparatus

TG–DTA analysis of the precursors was obtained using a thermogravimetric analyzer (model Netzsch STA 409) at a heating rate 0.1 °C/s under oxygen flow. Elemental analysis of the final products was determined using an induced coupled plasma (ICP) mass spectrometer (model VG Plasma Quad II-S) by dissolving the given powder samples in HCl. The phase identification was carried out by XRD using a Philips X’Pert PRO Materials Research Diffractometer (PW3050) equipped with a Cu anticathode (CuKα radiation λ = 1.54056 Å) at room temperature. The measurements were recorded under Bragg–Brentano geometry at 2θ with step 0.05° in the range 10–80°. The XRD data were analyzed by the Rietveld profile refinements using the Fullprof program. The particle morphology of the LiNi₀.5Mn₁.5O₄ powders was examined with a scanning electron microscope (SEM, Philips XL30). Magnetic measurements were performed in a SQUID magnetometer (Quantum Design MPMS-5S). Zero-field-cooled and field-cooled magnetic susceptibility data were obtained in a field at 1 T in the temperature range 4 ≤ T ≤ 300 K.

The FTIR absorption spectra were recorded with a Fourier transform interferometer (model Bruker IFS113v) in the wavenumber range 120–1,200 cm⁻¹ at a spectral resolution 2 cm⁻¹. Samples were ground to fine powders and dispersed into CsI pellet. RS spectra were collected with a Raman-laser spectrometer (Jobin–Yvon model U1000) equipped with holographic grating double-monochromator coupled with a spatial filter. Powders were press-compact into pellet to form a mirror-like surface illuminated with the 514.5-nm emission from an Ar-ion laser (model Spectra-Physics 2020). Each RS spectrum was the average of 10 scans collected at a spectral resolution of 1 cm⁻¹. The curve fitting is based on the original algorithm of nonlinear peak fitting described by Marquardt, and known as the Levenberg–Marquardt method [31], by using the GRAM/386 software from Galactic Industries. The fitting calculations were done assuming a linear baseline for the spectra and that all the Raman lines introduced in the fit had a mixed Gauss–Lorentz line shape [32].

The electrochemical properties of the product were tested at room temperature in cells with metallic lithium as anode electrode. Measurements were carried out following the experimental procedure previously described [33] using Teflon laboratory-cell hardware and the Mac-Pile system. Quasi open-circuit voltage curves were recorded using a current pulse of 0.1 mA/cm² applied for 1 h followed by a relaxation period of 0.5–1.0 h, which allows the recording of the transient voltage for the determination of chemical diffusion coefficients of Li⁺ ions in the LiNi₀.5Mn₁.5O₄ lattice.

Results and discussion

Synthesis, structure, and morphology

ICP measurements were carried out on LiNi₀.5Mn₁.5O₄ oxides for determination of their elemental analysis. As shown in Table 1, both samples have a composition near the ideal one. These results prove that (1) the wet-chemical

| Sample     | Li (wt.%) | Ni (wt.%) | Mn (wt.%) | Composition                  | SG       | a (Å)   |
|------------|-----------|-----------|-----------|------------------------------|----------|---------|
| Sol-gel    | 3.81      | 15.80     | 45.35     | LiNi₀.49Mn₁.49O₄             | Fd3m     | 8.182 (7) |
| Pyrolysis  | 3.83      | 15.70     | 45.15     | LiNi₀.49Mn₁.5O₄              | P4/122   | 8.168 (5) |
| Theoretical| 3.80      | 16.07     | 45.11     | LiNi₀.5Mn₁.5O₄               | –        | –       |

Fig. 2 TG and DTA of LiNi₀.5Mn₁.5O₄ grown by a sol-gel method via citric acid and b combustion method via glycine.
synthesis assisted by glycine provides better stoichiometric samples than succinic acid and (2) no significant loss of lithium oxide has occurred during the thermal treatment.

Figure 2a,b shows the TG–DTA curves of LiNi_{0.5}Mn_{1.5}O_{4} powders grown by the sol-gel method via citric acid and the pyrolysis synthesis via glycine, respectively. The thermal analyses of materials grown by wet-chemical methods are characterized by three steps: (1) the removal of water up to 200 °C, (2) a strong exothermic reaction in the range 240–300 °C, and (3) the region of crystallization of the final product LiNi_{0.5}Mn_{1.5}O_{4} above 400 °C.

In both methods, the strong exothermic peak appears at 295 (sol-gel) and at 245 °C (pyrolysis) after the departure of the remaining water from the xerogel. The exothermic heat generated from the decomposition of acetate-ions xerogel and the combustion of the carboxylic acid, i.e., citric acid or glycine, is predicted to be utilized as the lattice energy required for the formation of LiNi_{0.5}Mn_{1.5}O_{4} phase. More than half of the weight loss occurs during this stage because of a violent oxidation–decomposition reaction. It appeared that citric acid and glycine acted as a fuel in the calculation of the gel precursor, accelerating the decomposition of acetate ions. The gel precursor was self-burning once ignited because the decomposed acetate ions acted as an oxidizer. The crystallization starts below 400 °C; thus, a well-crystallized and pure LiNi_{0.5}Mn_{1.5}O_{4} phase was obtained at 680 °C. While the calcination process at this stage was very complicated, it could be presumed that the last weak exothermic at ca. 460 °C in the DTA curve corresponds to the crystallization of the LiNi_{0.5}Mn_{1.5}O_{4} phase. During this process the mass precursor darkened progressively as a result of oxidation.

Figure 3 presents the XRD patterns of LiNi_{0.5}Mn_{1.5}O_{4} phases grown by the citrate and the glycine route. The diagram of the normal spinel LiMn_{2}O_{4} (Fd3m SG) is shown for comparison. The LiNi_{0.5}Mn_{1.5}O_{4} powders exhibit phase-pure cubic spinel without impurities as far as XRD is concerned. When products are calcined above 400 °C, the diffraction patterns of LiNi_{0.5}Mn_{1.5}O_{4} grown by the sol-gel method can be indexed by comparing those of LiMn_{2}O_{4} in the cubic system with Fd3m symmetry. The results of structure refinements show a decrease of the lattice parameter which falls from $a=8.243$ Å for LiMn_{2}O_{4} to $a=8.182$ Å for LiNi_{0.5}Mn_{1.5}O_{4} sol-gel phase.

The diffraction patterns of LiNi_{0.5}Mn_{1.5}O_{4} grown by the pyrolysis method have been indexed in the cubic $P4_{1}32$ symmetry rather than the cubic Fd3m SG due to additional weak lines located at $2\theta = 15.3, 39.7, 45.7,$ and $57.5^\circ$, which are absent from that of the Fd3m structure, and the lowering symmetry observed in the resonance spectroscopy (Raman and FTIR), which could not be detected by XRD due to the small contrast of XRD patterns. Similar results have been reported for LiNi_{0.5}Mn_{1.5}O_{4} and LiMg_{0.5}Mn_{1.5}O_{4} [16, 34]. Therefore, the primitive unit cell in the cubic system with the $P4_{1}32$ (or $P4_{2}32$) SG results in the superstructure from the 1:3 cation ordering. The cubic cell parameter $a=8.1685$ Å is in agreement with previously reported values which describe an octahedrally ordered spinel structure [16–19]. We assume similar distribution of the atoms, namely, Li on 8c, Ni on 4b, Mn on 12d, O(1) on 24e, and O(2) on 8c Wyckoff positions. The SG $P4_{1}32$ allows placing the larger Ni^{2+} ions (ionic radius 0.69 Å) in the bigger 4b site instead of the 16d site of normal spinel structure. The smaller unit cell dimension is primarily due to the change in the Mn oxidation state. Despite the replacement of a fraction of Mn ions by bigger Ni^{2+} ions, the Mn valence change effect prevails [5]. Recently, Strobel et al. [6] suggested that the occurrence of cation ordering corresponds to the largest valence difference $\Delta Z$ (Mn–M)≥2. The net result is thus a significant optimization of space occupation, leading to a reduced unit cell volume.

In general, several factors can contribute to the broadening of peaks in X-ray diffraction [35]. For example, instrumental factors related to the resolution and the incident X-ray wavelength, as well as sample factors such as crystallite size and nonuniform microstrain, can cause line broadening. In the case of instrumental broadening, line width will vary smoothly with $2\theta$ or $d$ spacing. On the other hand, line broadening that originates from sample characteristics will have a different relationship. By combining the Scherrer’s equation for crystallite size and
the Bragg’s law for diffraction, crystallite size and microstrain components are estimated by using the following equation:

$$B^2 \cos^2 \theta = 16 \langle e^2 \rangle \sin^2 \theta + \frac{K^2 \lambda^2}{L^2},$$  \hspace{1cm} (1)

where $B$ is the full-width at half-maximum after correction of instrumental broadening for finely powdered silicon powder, $\theta$ is the diffraction angle, $\langle e^2 \rangle$ denotes local strains (defined as $\Delta d/d$ being the interplanar spacing), $L$ is the crystallite size, and $K$ is a near-unity constant related to crystallite shape. Plots of eight reflections for the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ and LiMn$_2$O$_4$ are presented in Fig. 4. The slope, $16 \langle e^2 \rangle$, and intercept, $K^2 \lambda^2/L^2$, were used to estimate the distortion and size parameters. Apparently, the integral breadth of Bragg lines of the doped spinel is slightly greater than those of pure samples, which suggests a smaller grain size and a higher microstrain content in LiNi$_{0.5}$Mn$_{1.5}$O$_4$ (Table 2).

The shape and morphology of the particles in spinel samples obtained by different methods were followed by SEM. Figure 5 shows the SEM images of the products calcined at 680 °C. It is shown that the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ powders consist of well-developed particles. The average particle size is estimated to be 200–280 nm for materials prepared by the sol-gel technique and 300–400 nm for materials prepared by the pyrolysis method. Moreover, the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ powders exhibit a uniform particle size distribution.

Raman and FTIR studies

Structural properties of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ phases grown by wet-chemical route with replacement of exactly $1/4$ B site (octahedral cations) were also investigated by vibrational spectroscopy. Because Raman and FTIR are local probes sensitive to the crystal symmetry, they are useful tools when poor chemical contrast of XRD prevents the determination of the structure (amorphous structure, cation ordering, etc.). Figure 6 shows the RS spectra of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ powders displaying obvious differences to their LiMn$_2$O$_4$ counterparts. The RS spectrum of LiMn$_2$O$_4$ consists of a series of broad bands between 300 and 700 cm$^{-1}$. The peaks at 625, 580, 483, 426, and 362 cm$^{-1}$ are assigned to $A_{1g}$, $F_{2g}^{(3)}$, $F_{2g}^{(2)}$, $E_g$, and $F_{2g}^{(1)}$ modes, respectively, as predicted by group theory for a cubic compound [36].

The introduction of Ni$^{2+}$ ions in the spinel lattice ($Fd3m$ SG) has modified the Raman spectra in a complicated manner. The following observations can be made: (1) The 625-cm$^{-1}$ band that is associated with the symmetric Mn-O stretching vibration of MnO$_6$ groups, 580, 483, 426, and 362 cm$^{-1}$ are assigned to $A_{1g}$, $F_{2g}^{(3)}$, $F_{2g}^{(2)}$, $E_g$, and $F_{2g}^{(1)}$ modes, respectively, as predicted by group theory for a cubic compound [36].

Table 2 Structural parameters of LiMn$_2$O$_4$ and LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinels prepared by wet-chemical method from the citrate precursor route

| Sample       | $a$ (Å) | $L$ (nm) | $\langle e^2 \rangle \times 10^6$ |
|--------------|--------|---------|-------------------------------|
| LiMn$_2$O$_4$| 8.224 (0) | 320     | 1.53                          |
| LiNi$_{0.5}$Mn$_{1.5}$O$_4$ | 8.182 (7) | 255     | 1.67                          |

Fig. 5 SEM images of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ phases grown by wet-chemical route via citrate precursors. a Normal spinel ($Fd3m$) and b ordered spinel ($P4_{1}32$). We note the difference in magnification between the images.
the decrease in the unit cell volume. The small shift of the symmetric stretching vibration of MnO$_6$ groups can be viewed as the shortness of Mn-O bond lengths and polyhedral distortion occurring in LiNi$_{0.5}$Mn$_{1.5}$O$_4$. The intensity of the shoulder located at 580 cm$^{-1}$ is enhanced upon nickel substitution. This may be due to the change of the Mn$^{3+}$/Mn$^{4+}$ ratio vs Ni$^{2+}$ in the material. The peak splitting between bands at 426 and 382 cm$^{-1}$ is reduced between the bands at 407 and 378 cm$^{-1}$ for LiNi$_{0.5}$Mn$_{1.5}$O$_4$ due to the polyhedral distortion.

For well-controlled synthesis of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ grown by glycine-assisted pyrolysis method, RS spectrum indicates characteristic peaks originating from the formation of a superlattice. Analytical results are in accordance with the P4$_i$32 SG in which Ni$^{2+}$ ions are located at the 4a sites in the cubic-close packed oxygen array. Comparison with the disordered LiNi$_{0.5}$Mn$_{1.5}$O$_4$ (Fd3m SG) compound is made in Fig. 6. The main driving force for octahedral cation ordering is believed to be the charge difference between Mn and Ni atoms. For Ni substitution, the RS patterns show a dramatic increase in the number of active modes. Considering the symmetry correlation $O^h_7 \rightarrow O^v$, the irreducible representations for normal spinel ($\Gamma_{NS}$) and for ordered spinel ($\Gamma_{OS}$) are given by (Table 3).

\[
\Gamma_{NS} = A_g(R) + E_g(R) + 3F_{2g}(R) + 4F_{1u}(ir), \quad (2)
\]

\[
\Gamma_{OS} = 6A_1(R) + 14E(R) + 20F_1(\text{ir}) + 22F_2(R). \quad (3)
\]

It is obvious that the number of expected Raman-active modes ($6A_1 + 14E + 22F_2$) is larger than for LiMn$_2$O$_4$ ($A_g + E_g + 3F_{2g}$). This prediction is well verified experimentally. The LiNi$_{0.5}$Mn$_{1.5}$O$_4$ phase clearly shows an ordering of octahedral cations into 4b and 12d sites in SG P4$_i$32 ($O^v$ spectroscopic symmetry). Both contain small fractions of lithium on the 4b site [6]. The sharpness of the Raman bands of the LiMn$_{3/2}$Ni$_{1/2}$O$_4$ sample is the fingerprint of well separated Ni and Mn sites resulting from the symmetry lowering. As the integer valence distribution is (Li$^+$(Ni$^{2+}$)$_{0.5}$(Mn$^{4+}$)$_{1.5}$O$_4$, the rather large broadness of the high-wavenumber band in normal spinel LiMn$_2$O$_4$ becomes a well-resolved triplet in ordered spinel lattice. This vibration is recorded at 500 cm$^{-1}$ in the case of NiO [37]. Figure 7 shows the FTIR absorbance spectra of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ powders compared with that of undoped LiMn$_2$O$_4$ normal-spinel. Frequencies have been determined by fitting using a mixed Gauss–Lorentz line shape. The FTIR spectrum of LiMn$_2$O$_4$ consists of two broad bands $\nu_1$ and $\nu_2$ at ca. 615 and 513 cm$^{-1}$. Weak shoulders are observed on the low-wavenumber side with the

Table 3 The irreducible representation of the vibrational active modes of the 1:3-ordered spinel phase in $O^v$ spectroscopic group

| Atom | Wyckoff position | Point group | Irreducible representation (point group $\rightarrow O^v$ correlation) |
|------|-----------------|-------------|-----------------------------------------------------------|
| Li   | 8c              | C$_3$       | $A_1$  \rightarrow $A_1 + 2E + 3F_{1y} + 3F_{2y}$ |
| Mn   | 12d             | C$_2$       | $A_1$  \rightarrow $A_1 + 2E + 5F_{2y} + 5F_{2z}$ |
| Ni   | 4b              | D$_3$       | $A_1$  \rightarrow $A_1 + 2E + 3F_{1y} + 3F_{2y}$ |
| O1   | 8c              | C$_i$       | $A_1$  \rightarrow $A_1 + 2E + 3F_{1y} + 3F_{2y}$ |
| O2   | 24e             | C$_i$       | $A_1$  \rightarrow $A_1 + 2E + 3F_{1y} + 3F_{2y}$ |
| Total| 6A$_1$ + 7A$_2$ + 14E + 23F$_{1y}$ + 22F$_{2y}$ | 3F$_{1y}$ |
| Acoustic | 3F$_{1y}$ | | |
| Inactive | 7A$_2$ | | |
| Raman (42 modes) | 6A$_1$ + 14E + 22F$_{2y}$ | | |
| Infrared (20 modes) | 20F$_{1y}$ | | |

Fig. 6 RS spectra of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ powders compared with that of undoped normal-spinel (Fd3m). The 1:3 ordered phase (P4$_i$32) gives a very different spectrum with much sharper lines resulting from the symmetry lowering. Spectra were fitted using the curve analysis based on a nonlinear peak fitting known as the Levenberg–Marquardt method with Lorentzian line shape.
component $\nu_3$ at 420 cm$^{-1}$. Four infrared active modes ($F_{1u}$) are allowed in the $O_7^h$ spectroscopic symmetry [36]. These results displayed in Fig. 7 show that the local structure of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ is strongly affected by the presence of Ni$^{2+}$ ions. The FTIR spectrum of the sample grown by glycine-assisted method presents nine well-resolved modes as predicted by the symmetry lowering $O_7^h \rightarrow O^7$. An increase in the number of infrared-active modes confirms the previous RS data.

Magnetic studies

Figure 8 shows the temperature dependence of the reciprocal molar magnetization normalized by magnetic fields ($H/M$) under an applied field $H=10$ kOe for the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinel grown by wet-chemical route. A ferromagnetic ordering is clearly shown at the Curie temperature $\theta=129$ K for Ni-doped spinel. The magnetization was measured at 10 kOe gradually at lower temperatures on field cooling and measurements on heating after zero-field cooling.

Figure 9 displays the field dependence of the magnetization for the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ sample at different fixed temperatures. Below 150 K, nonlinearity is clearly observed, suggesting the presence of a strong ferromagnetic contribution. At low temperatures, a typical magnetization curve shows the rapid increase with the magnetic field; the magnetic saturation is obtained at a relatively small field. Several authors have reported such a magnetic behavior. Blasse [38] has shown the ferromagnetic spin alignment below 130 K in LiNi$_{0.5}$Mn$_{1.5}$O$_4$. Strobel et al. [5] found a large increase in the magnetic susceptibility at low temperature, starting at ca. 120 K on cooling. These features were attributed to the formation of clusters resulting from a strong magnetic frustration in the structure. Recent results by Nakamura et al. [39] have shown that the asymptotic Curie temperature moved in the positive direction and changed its sign from negative to positive around $y=0.2$ in LiNi$_y$Mn$_{2-y}$O$_4$ spinels. For LiNi$_{0.5}$Mn$_{1.5}$O$_4$, the ferromagnetism is reported with a saturation magnetization of 105 emu/g at 5 K.

Due to a large ferromagnetic component, the magnetization curves, $M(H)$, are not linear functions of the applied
magnetic field, as shown in Fig. 9. The first consequence is an ambiguity in what is called magnetic susceptibility, $\chi_m$, because $M/H$ is distinct from $dM/dH$. Let us consider two options for the ferromagnetic behavior of LiNi$_{0.5}$Mn$_{1.5}$O$_4$. First, the material should be considered as a ferromagnet itself, i.e., intrinsic ferromagnetism, due to the couplings Mn$^{4+}$-$\text{O}^{2-}$-Mn$^{4+}$ and Ni$^{2+}$-$\text{O}^{2-}$-Ni$^{2+}$. Secondly, ferromagnetic particles could be included in the spinel framework; such impurities (Mn$_3$O$_4$) in small concentrations that cannot be detected by XRD. In this case, one has an extrinsic ferromagnetism. If such particles are present in LiNi$_{0.5}$Mn$_{1.5}$O$_4$, the ferromagnetism originates from these “clusters.” Ferrimagnetic impurities such as NiO or Li$_{z}$Ni$_{1-z}$O are excluded because they possess very high Curie temperatures, e.g., $T_c = 2,000$ K [38]. In this case, the Curie–Weiss law is valid only at $T > |\theta|$. This condition can only be fulfilled above room temperature for LiNi$_{0.5}$Mn$_{1.5}$O$_4$, which precludes a quantitative analysis of the $H/M$ curve. The Curie temperature is determined by experiments (see Fig. 8). At the lowest temperature $T=4$ K, the magnetization saturates, so that

$$M(T=4K) = Nn\mu_{eff}(0),$$ \hspace{1cm} (4)

where $n$ is the number of magnetic ions of effective moment $\mu_{eff}$. The magnetic moment is equal to 3.02 $\mu_B$, corresponding to a specific magnetization 88 emu/g for LiNi$_{0.5}$Mn$_{1.5}$O$_4$. This value is consistent with the data reported by Nakamura et al. [39].

Figure 10 shows the ESR spectra of LiMn$_2$O$_4$ and LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinels measured at room temperature. The ESR spectrum of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ shows a much narrower signal ($\Delta H_{pp}=170$ mT) centered at $g=2.0$. The band has a complex shape with two components. The signal is attributed to Mn$^{4+}$ ions, which are the only paramagnetic entities in this compound.

**Electrochemical studies**

The electrochemical properties of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinels have been widely reported [13–25]. Also, in this section, we briefly present the electrochemical data of our samples prepared by wet-chemical methods. The first charge–discharge profiles of the spinel Li/LiNi$_{0.5}$Mn$_{1.5}$O$_4$ cells in the voltage range 3.5–4.9 V vs Li$^0$/Li$^+$ have been recorded for normal $A[B_2]O_4$ spinel lattice (Fig. 11a) and for ordered 1:3 spinel framework (Fig. 11b). Our data are consistent with those of prior results in terms of their high-voltage characteristics [41]. As can be seen in Fig. 11, the LiNi$_{0.5}$Mn$_{1.5}$O$_4$ samples reveal an operating voltage higher
than 4.5 V. They reveal an obvious two-step lithium intercalation/deintercalation behavior. A small plateau appears at ca. 4.0 V, related to the redox couple Mn$^{3+}$/Mn$^{4+}$. The appearance of the voltage plateau at 4.7 V has been attributed to the oxidation of Ni$^{2+}$ to Ni$^{4+}$ [4]. The normal-spinel phase shows predominantly one-step reaction at 4.65 V, while the voltage profile of the ordered-spinel structure transforms from a sloping curve to a flat curve at 4.72 V. In the voltage range 3.5–4.9 V, the Li//LiNi$_{0.5}$Mn$_{1.5}$O$_4$ delivers a capacity 133 mAh/g during the first discharge.

Figure 12 displays the plot of the incremental derivative capacity ($-\partial Q/\partial V$) for the Li//LiNi$_{0.5}$Mn$_{1.5}$O$_4$ cell as a function of the cell voltage vs Li$^0$/Li$^+$ for the first charge–discharge process.

In this work, we have developed wet syntheses such as sol-gel and pyrolysis assisted by carboxylic acid to grow LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinel materials. Structural and electrochemical characterizations were made and properties were compared with LiMn$_2$O$_4$ spinel. The overall picture that emerges from these investigations is as follows.

Depending of the preparation procedure, two varieties of samples were obtained: the normal-spinel structure (via citric acid) and the ordered-spinel structure (via glycine). These structures refined with the Fd3m and P4_132 SGs, respectively, show a smaller cubic unit cell than the undoped LiMn$_2$O$_4$ spinel. The shortening of M-O bonds is observed in Raman spectra by a shift of the stretching mode of MnO$_6$ entities. The superstructure of LiNi$_{0.5}$Mn$_{1.5}$O$_4$ material grown by pyrolysis method is confirmed by local probes, such as Raman and FTIR.

Raman measurements reveal that the oxidation state of Ni was +2 in LiNi$_{0.5}$Mn$_{1.5}$O$_4$ spinels. Magnetic measurements have evidenced the ferromagnetic ordering below $T_c$=129 K in LiNi$_{0.5}$Mn$_{1.5}$O$_4$. The replacement of the 3d$^8$-Ni for Mn in LiNi$_{0.5}$Mn$_{1.5}$O$_4$ oxides is beneficial for its charge–discharge cycling performance. For a cut-off voltage 3.5–4.9 V, the electrochemical capacity of the Li//LiNi$_{0.5}$Mn$_{1.5}$O$_4$ cell is ca. 133 mAh/g during the first discharge.

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