1. Introduction

The interaction between prey species and predator species constitutes biological processes of major ecological importance that are ubiquitous in nature. The dynamical interaction between prey-predator models with different kinds of response function is a dominant theme in applied mathematics and theoretical ecology [1]. Mathematical modeling for prey-predator interplays has drawn attention to the mathematicians since the pioneering worked by Lotka and Volterra in the year 1920s, and there has been a large investigation for their rich dynamics [2–6]. When we go through the interaction between prey-predator system through mathematical modeling, initially it appears to be very straightforward, but at the end, it becomes very challenging tasks from the view point of mathematicians. The investigation and validation for the proper ecological behavior of the proposed model under consideration becomes most challenging and crucial. Since the complicated and rich dynamics for the interactive populations are usual throughout the world, many ecologists have studied the processes that influence the kinetics of predator-prey system and are interested to know which model would be best to represent the interaction among species.

The dynamical behavior for the prey-predator system takes part in a significant role in theoretical ecology. There are many ecologists feeling that the unique nonnegative steady state for a predator-prey model is globally asymptotically stable if it becomes stable asymptotically although it is not always true. It can be shown that, under suitable condition, a unique nonnegative locally asymptotically stable steady state has at least one limit cycle around the steady state. Therefore, many mathematicians attempt to utilize some well-studied methods to obtain situations of global stability of the steady state for the prey-predator model [5–10].

Generally, the prey-predator system obeys two basic principles: (i) prey species has natural birth rate and both the...
species has natural death rate; (ii) predator species can only grow by consuming the prey species [11]. In order to explain the predation phenomenon, Holling [12] has proposed different types of functional responses such as type-I, type-II, and type-III. The response functions are the rate of prey consumption by the predator per unit time. The functional responses proposed by Holling [12] are the functions of prey density only. The response functions are independent of predator density. One of the widely studied functional responses is Holling type-II response function, which is characterized by decelerating intake rate. The rate of predation rises as prey density increases, but after a certain stage, the rate of predation remains constant although prey density increases. In the nature, predator population death rate is likely to be maximum when the prey population is minimum due to the lack of food. Later, it is recognized that the response function is also dependent on predator density. In the year 1975, Beddington and DeAngelis [13, 14] introduce a response function, which is alike to Holling type-II response function but holds an additional term, which represents the mutual interaction between the predators. Hence, Beddington–DeAngelis-type response function is a function of both prey and predator density; that is, the rate of predation is function of both prey and predator density. However, depending upon the structural complexity of the prey habitat, the response function for the prey-predator interaction may be different [15].

Recently, a series of mathematical models has been developed to better understand the dynamics of the predator-prey system with different kinds of functional responses [6, 16–24]. Dalziel et al. [16] studied a modified Holling type-II predator-prey model, based on the premise that the search rate of predators is dependent on the prey density, rather than constant. They performed a complete analysis of the global behavior of their model and the dichotomy properties. Wang et al. [17] investigated an improved predator-prey model that incorporates seasonality and a predator maturation delay, leading to a system of periodic differential equations with a time delay. They also defined the basic reproduction number $R_0$ and show that it is a threshold parameter determining whether the predators can coexist with the prey species. Kohnke et al. [18] studied a predator-prey model and developed a functional response using timescale separation. In their study, they showed that the resulting functional response contains a single parameter that controls whether the group defense functional response is saturating or dome-shaped. Their model exhibits bifurcation analysis. Antwi-Fordjour et al. [19] investigated the interplay of the mutual interference exponent, and prey refuge, on the behavior of a predator-prey model with a generalized Holling-type functional response considering in particular the “nonsmooth” case. The authors studied the dynamical properties of their model and derived conditions for the occurrence of saddle-node, transcritical, and Hopf bifurcations. Also, they derived the sufficient condition for finite time extinction of the prey population.

When we look through literature review for the published research papers on the dynamics of the prey-predator model, we notice that there are multiple numbers of good research works on the dynamics of prey and predators with different kinds of functional responses. Some researchers investigated and elevated some open quires of various kinds of response function. Seo and Kot [2] have compared two prey-predator models with Holling type-I response function. In [25], the author investigated a predator-prey model with different kinds of functional responses by incorporating the role of constant prey refuge. Aiello and Freedman [26] studied a single species model with the effect of discrete time lag, in which the species was divided into mature and immature. Their model demonstrates nonnegative equilibria as the global attractor. Kon et al. [27] established the criteria for the permanence of stage structure prey-predator interactions. These are some existing articles which focused on Beddington–DeAngelis-type functional responses [28–31]. Beddington–DeAngelis-type response function confesses plentiful and biologically relevant prey-predator dynamics [29], which attract researchers to further study the predator-prey model including Beddington–DeAngelis response function. Khajanchi [28] studied a predator-prey model with age-structured incorporating Beddington–DeAngelis-type response function, in which the author performed global stability analysis. Xie et al. [20] studied a predator-prey system by using a system of fractional order differential equations with Holling type-III response function and incorporate the effect of discontinuous harvesting. Predator-prey system has been investigated by introducing discrete time delay with disease in the predator by Huang et al. [21]. Also, nonautonomous Leslie–Gower predator-prey system with nonlinear prey harvesting was studied by Alam et al. [22]. Sarkar and Khajanchi [23] studied a predator-prey system by incorporating the effect of fear with Holling type-II response function, where the authors preformed the bifurcation analysis of the system to better understand the dynamics of the model. Seasonally, varying prey-predator system with the Allee effect model was investigated by Rebelo and Soresina [24], and they studied the coexistence of all the species. Time delay plays an important role in stability switching and may change drastically the qualitative behavior of solutions for the model system under consideration [32–35].

The main aim of this manuscript is to study a very simple two-tier continuous predator-prey model with different kinds of functional responses. Although the model is very simple, it gives rich and biologically meaningful dynamics. A comparative study of the stability property was done by considering Holling type-I, Holling type-II, and Beddington–DeAngelis-type response functions. Our predator-prey system experiences Hopf bifurcation with respect to the prey birth rate $r_0$, which gives more complicated dynamics of the system. The predator-prey system also experiences transcritical bifurcation with respect to the prey birth rate $r_0$. The biological implications of the theoretical and numerical results are also investigated.

This article is organized in the following way: we present a predator-prey model with different kinds of functional responses in the next section. In Section 3, we discuss the qualitative behavior for the model, and positivity for the model, boundedness for the solutions, and the position of all
possible feasible equilibria and the existence criteria of the equilibria are presented. The criteria for stability of the predator-prey system around these equilibria are studied. In the same section, we performed a comparative study by considering different kinds of functional responses. We performed global stability analysis by constructing suitable Lyapunov function. By using normal form theory and center manifold theorem, we investigate direction and stability of Hopf bifurcation in Section 4. In the next section, we give some numerical simulations to support our analytical findings. Section 6 presents the discussion about key findings for our problem.

2. The Model

The interaction between prey and predator species in simplest form is represented by the following system of coupled differential equations [2, 36]:

\[
\begin{align*}
\frac{du}{dt} &= uf(u) - g(u, v)v, \\
\frac{dv}{dt} &= v(-\delta_2 + \theta g(u, v)),
\end{align*}
\]

where \(v(t)\) and \(u(t)\) represent the predator and prey populations, respectively, at any time \(t\) and \(f(u)\) stands the birth rate of prey species in exclusion of any predator population as well as death rate and the intraspecific competition between the prey species. Here, \(\delta_2\) denotes the natural death rate for predator species and \(\theta\) stands the conversion coefficient of prey’s biomass to predator’s biomass, and \(g(u, v)\) indicates the prey-dependent response function.

Our proposed model consists of two population, namely, prey population and predator population. Concentration of prey and predator species at time \(t\) can be represented by \(u(t)\) and \(v(t)\), respectively. Both the prey and predator populations are continuous in time \(t\). The prey-predator interaction with response function \(g(u, v)\) can be expressed by the system of coupled differential equations:

\[
\begin{align*}
\frac{du}{dt} &= r_0u - \delta_1u - \gamma u^2 - g(u, v)v, \\
\frac{dv}{dt} &= v(-\delta_2 + \theta g(u, v)).
\end{align*}
\]

With initial values \(u(0) = u_0 \geq 0\) and \(v(0) = v_0 \geq 0\), all the model parameters \(r_0, \delta_1, \gamma, \delta_2,\) and \(\theta\) are positive constants and can be interpreted as follows: \(r_0\) represents the birth rate of prey species in absence of predators, \(\delta_1\) represents the death rate for prey population, \(\gamma\) denotes the intraspecific competition of prey species, \(\delta_2\) represents the death rate for predator population, and \(\theta\) represents the conversion coefficient of prey’s biomass to predator’s biomass. The dimension of prey and predator biomass is Mass. Also, the dimensions of \(r_0, \delta_1, \gamma, \) and \(\delta_2\) are \(\text{Time}^{-1}\), \(\text{Time}^{-1}\), \(\text{Mass}^{-1}\), \(\text{Time}^{-1}\), and \(\text{Time}^{-1}\), respectively, and the parameter \(\theta\) is dimensionless. In exclusion of predator species, the prey population obeys the logistic growth with an intrinsic growth rate \((r_0 - \delta_1) > 0\).

2.1. Functional Response. In ecology, the functional responses are the intake rate of predator population as a function of prey population. Holling [12] proposed three types of response function, namely, Holling types-I, II, and III. In our model formulation, we consider Holling Type-II response function in the following form:

\[
g(u, v) = \frac{\alpha u}{1 + \beta u},
\]

where \(\alpha\) represents the attack rate and \(\beta\) is the handling time to capture the prey population; both \(\alpha\) and \(\beta\) are positive constants. The dimensions of \(\alpha\) and \(\beta\) are \(\text{Mass}^{-1}\text{Time}^{-1}\) and \(\text{Mass}^{-1}\), respectively. The predation rate increases as prey density increases, but after a certain stage, rate of predation becomes constant although prey density increases. The functional response depends on prey density but not on predator density. By using the functional form of \(g(u, v)\) given by (3), system (2) becomes

\[
\begin{align*}
\frac{du}{dt} &= r_0u - \delta_1u - \gamma u^2 - \frac{\alpha uv}{1 + \beta u} = F_1(u, v), \\
\frac{dv}{dt} &= \frac{\alpha uv}{1 + \beta u} - \delta_2v = F_2(u, v).
\end{align*}
\]

3. Qualitative Behavior of the Model

3.1. Positivity, Boundedness, and Uniform Persistence. Right side of the system (4) is a continuous function for dependent variable \(t\). Hence, after integration on both sides of each equation of system (4), we get

\[
\begin{align*}
\frac{du}{dt} &= u(0)\exp\left(\int_0^t \left[ r_0 - \delta_1 - \gamma u(s) - \frac{\alpha v(s)}{1 + \beta u(s)} \right] ds \right) > 0, \quad \forall t \geq 0, \\
\frac{dv}{dt} &= v(0)\exp\left(\int_0^t \frac{\alpha uv(s)}{1 + \beta u(s)} - \delta_2 \right) ds > 0, \quad \forall t \geq 0.
\end{align*}
\]

From the above expressions of \(u(t)\) and \(v(t)\), it is clear that both \(u(t)\) and \(v(t)\) remain nonnegative for all finite time, that is, \(t \geq 0\). Hence, \(\mathbb{R}_+^2 = \{(u(t), v(t)) : u(t) \geq 0, v(t) \geq 0\}\) is positively invariant of system (4). In order to prove the boundedness of the prey-predator model (4), we consider the function
\[ \phi(t) = u(t) + \frac{1}{\theta} v(t). \]  

(6)

Differentiation with respect to \( t \) gives

\[ \frac{d\phi(t)}{dt} = \frac{du(t)}{dt} + \frac{1}{\theta} \frac{dv(t)}{dt} = (r_0 - \delta_1)u(t) - \gamma u^2 - \frac{\delta_2}{\theta} v(t). \]

(7)

Now, for each \( \phi(t) > 0 \), we get

\[ \phi(t) + \rho \phi(t) \leq r_0 u - \gamma u^2, \]

where \( \rho = \min\{\delta_1, (\delta_2/\theta)\} \). The maximum value of \( r_0 u - \gamma u^2 \) is \( (r_0^2/4\gamma) \). Therefore,

\[ \frac{\phi(t)}{dt} + \rho \phi(t) \leq \frac{r_0^2}{4\gamma} = m \text{ (say)}. \]

(9)

Due to the theory of differential inequality, we have

\[ 0 \leq \phi(t) \leq \frac{m}{\rho} (1 - e^{-\rho t}) + \phi(0) e^{-\rho t} \]

(10)

for \( t \rightarrow \infty \), and we have \( 0 \leq \phi(t) \leq (m/\rho) \). Therefore, \( \phi(t) \) is bounded. Thus, all the solution of (4) are confined in

\[ \Omega = \left\{ (u, v) : 0 \leq u(t) + \frac{1}{\theta} v(t) \leq \frac{m}{\rho} + \varepsilon, \quad \forall \varepsilon \geq 0 \right\}. \]

(11)

This implies that the solution of system (4) is bounded. Thus, we have the following theorem.

**Theorem 1.** For any initial value \((u_0, v_0) \in \mathbb{R}_+^2\), all the solutions for (4) are nonnegative and bounded.

From the ecological perspective, the positivity of the prey-predator system (4) interprets that prey and the predator populations are biologically well-behaved. The boundedness for the solutions of the prey-predator model (4) implies that neither the prey population nor the predator population will grow unboundedly or abruptly for a large period.

**Definition 1.** System (4) is said to be uniformly persistent if there exist constants \( \kappa \) and \( \tau \), where \( 0 \leq \kappa \leq \tau \) such that

\[ \max \{\limsup_{t \to \infty} u(t), \limsup_{t \to \infty} v(t)\} \leq \tau \quad \text{and} \quad \min \{\liminf_{t \to \infty} u(t), \liminf_{t \to \infty} v(t)\} \geq \kappa, \]

where \( \kappa \) and \( \tau \) are independent of the initial conditions.

**Theorem 2.** System (4) is uniformly persistent if \( r_0 > \delta_1 + a\overline{\nu} \), where \( \overline{\nu} \) is the upper bound of \( \nu(\tau) \).

**Proof.** From the proof of Theorem 1, we have \( \limsup_{t \to \infty} u(t) \leq (m/\rho) \) and \( \limsup_{t \to \infty} v(t) \leq (m\theta/\rho) \). We choose \( \tau = \max\{(m/\rho), (m\theta/\rho)\} \). Hence, there exists \( \tau > 0 \) such that \( \max \{\limsup_{t \to \infty} u(t), \limsup_{t \to \infty} v(t)\} \leq \tau \).

Now, from the first equation of system (4), we have

\[ \frac{du}{dt} = r_0 u - \delta_1 u - \gamma u^2 - \frac{au}{1 + \beta u} \geq r_0 u - \delta_1 u - \gamma u^2 - a\overline{\nu}. \]

(12)

where \( \overline{\nu} \) is the upper limit of \( \nu \) in \( \Omega \).

Thus, we can conclude that \( \lim \inf_{t \to \infty} u(t) \geq (r_0 - \delta_1 - a\overline{\nu}/\gamma) \) with \( r_0 > \delta_1 + a\overline{\nu} \). Again, from the second equation of (4), we have

\[ \frac{dv}{dt} = a\overline{\nu} \frac{u}{1 + \beta u} - \delta_2 v \geq -\delta_2 v. \]

(13)

Hence, we have \( \lim \inf_{t \to \infty} v(t) \geq 0 \). Thus, if we choose \( \kappa = \min\{0, ((r_0 - \delta_1 - a\overline{\nu})/\gamma)\} \) with \( r_0 > \delta_1 + a\overline{\nu} \), then \( \min \{ \lim \inf_{t \to \infty} u(t), \lim \inf_{t \to \infty} v(t) \} \geq \kappa \). Therefore, system (4) will be uniformly persistent if \( r_0 > \delta_1 + a\overline{\nu} \). \( \square \)

### 3.2. Equilibria and Their Existence

The singular points for the model (4) are nonnegative solutions for \( F_1(u, v) = 0 \) and \( F_2(u, v) = 0 \) in \( \mathbb{R}_+^2 \); that is, the equilibrium points of system (4) are the positive point of intersection for zero growth of prey and predator isoclines. The prey isoclines are given by

\[ \alpha \theta \]

where \( \alpha = \min\{0, ((r_0 - \delta_1 - a\overline{\nu})/\gamma)\} \) with \( r_0 > \delta_1 + a\overline{\nu} \), and \( \beta = \alpha \theta \). System (4) possesses three biologically meaningful singular points, as follows:

(i) Trivial steady state \( E_0(0,0) \).

(ii) Boundary singular point \( E_1(\pi,0) \), where \( \pi = (r_0 - \delta_1)/\gamma \), which is feasible only when \( r_0 > \delta_1 \); that is, growth of prey population is higher than the death rate for prey population.

(iii) The positive interior equilibrium point \( E^*(u^*, v^*) \), where \( u^* = \delta_1/(\alpha \theta - \beta \delta_2) \) and \( v^* = (r_0 - \delta_1 - \gamma u^*) ((1/\alpha) + (\beta \alpha u^*) \).

The shape and position of the curve \( v = (r_0 - \delta_1 - \gamma u) (1 + \beta u)/\alpha \) depends upon multiple parameters. For different values of \( \beta \), the zero growth prey isoclines are convex curve joining the boundary equilibrium \( E_1(\pi,0) \) and \( (0, ((r_0 - \delta_1)/\alpha)) \) (see Figure 1(a)). As \( \beta \) increases, the vertical predator isocline shifted towards decreasing \( u \). For different values of \( \alpha \), zero growth prey isocline becomes a convex curve, which meets at boundary equilibrium \( E_1(\pi,0) \) (see Figure 1(b)). As \( \alpha \) increases, the vertical predator isocline shifted towards decreasing \( u \) and the prey isocline shifted towards decreasing \( v \). Also, for different values of \( r_0 \), the zero growth prey isoclines are convex curve, but the zero growth predator isoclines do not change position. As \( r_0 \) increases, the zero growth prey isocline shifted towards \( v \) increasing. In all these situations, the interior equilibrium point \( E^*(u^*, v^*) \) exists only if \( \alpha \theta - \beta \delta_2 > 0 \) (see Figure 1(c)).

### 3.3. Local Stability Analysis

In this section, we analyzed the linear stability for the predator-prey system (4). Using linearization techniques, the local behavior of the nonlinear predator-prey system can be investigated. Generally, the
The predator-prey system is linearized around every biologically feasible singular point, and the model is perturbed by a small quantity and observed whether the model returns to that state of singular point converges to another state of singular point. Analysis of local stability leads to investigate qualitative dynamics for the model system under consideration.

For local stability analysis of the considered model (4), first calculate the Jacobian matrix obtained from each of the equilibrium point. The variational matrix of the model (4) at \((u, v)\) is

\[
J = \begin{bmatrix}
    F_u & F_v \\
    G_u & G_v
\end{bmatrix},
\]

where

\[
F_u = r_0 - \delta_1 - 2\gamma u - \frac{\alpha v}{(1 + \beta u)^2},
\]

\[
F_v = \frac{\alpha u}{1 + \beta u},
\]

\[
G_u = \frac{\alpha \theta v}{(1 + \beta u)^2},
\]

\[
G_v = \frac{\alpha \theta u - \delta_2}{1 + \beta u}.
\]

**Theorem 3.** If \(r_0 < \delta_1\), then the trivial singular point \(E_0(0, 0)\) of (4) is locally asymptotically stable and otherwise unstable.

**Proof.** The eigenvalues of the Jacobian matrix (14) of the system (4) at the trivial equilibrium point \(E_0(0, 0)\) are \(r_0 - \delta_1\) and \(-\delta_2\). Here, \(\delta_2 > 0\); therefore, the trivial equilibrium point \(E_0(0, 0)\) of the model (4) will be locally asymptotically stable if \(r_0 < \delta_1\); that is, the growth rate of prey species is less than the natural death rate of prey species. It can be noted that if the trivial equilibrium point \(E_0(0, 0)\) is asymptotically stable, then the boundary singular point \(E_1(\bar{u}, 0)\) does not exist. For \(r_0 = \delta_1\), the trivial equilibrium point \(E_0(0, 0)\) has eigenvalues 0 and \(-\delta_2\). So \(E_0(0, 0)\) is a nonhyberbolic equilibrium point [6].

**Theorem 4.** If \(\delta_1 < r_0 < \delta_1 + \gamma \delta_2 / (\alpha \theta - \beta \delta_2)\), then the boundary equilibrium point \(E_1(\bar{u}, 0)\) of the model (4) is locally asymptotically stable and otherwise unstable.

**Proof.** The eigenvalues of the Jacobian matrix (14) of the system (4) at the equilibrium point \(E_1(\bar{u}, 0)\) are \(\lambda_1 = -(r_0 - \delta_1)\) and \(\lambda_2 = ((\alpha \theta (r_0 - \delta_1)) / (\gamma + \beta (r_0 - \delta_1))) - \delta_2\). The boundary equilibrium point \(E_1(\bar{u}, 0)\) will be locally

![Figure 1](image-url)
asymptotically stable if \( \lambda_1 < 0 \) and \( \lambda_2 < 0 \), and this condition implies \( \delta_1 < r_0 - \delta_1 + (\gamma \delta_2 / a \theta - \theta \beta \delta_2) \). Biologically, \( r_0 > \delta_1 \) implies that the rate of growth of prey population is higher than the natural rate of death of the prey. Also, \( r_0 < \delta_1 + (\gamma \delta_2 / a \theta - \theta \beta \delta_2) \) implies that the prey growth rate cannot be unbounded. It is to be noted that if the boundary equilibrium state \( E \) is locally asymptotically stable, then the trivial equilibrium state \( E_0 \) is unstable.

**Theorem 5.** If \( r_0 - \delta_1 - 2\gamma \alpha^* - (\alpha v^*/(1 + \beta u^*))^2 < 0 \), then the positive interior singular point \( E^* (u^*, v^*) \) of (4) will be asymptotically stable and otherwise unstable.

**Proof.** The variational matrix around the coexisting singular point \( E^* (u^*, v^*) \) for the system (4) is given by

\[
J^* = \begin{bmatrix}
 r_0 - \delta_1 - 2\gamma \alpha^* - \frac{\alpha v^*}{(1 + \beta u^*)^2} - \frac{a\alpha^*}{1 + \beta u^*} & a\alpha^* \\
\frac{a\alpha^*}{1 + \beta u^*}^2 & 0
\end{bmatrix}
\]

The characteristic equation of the above matrix \( J^* \) can be expressed as

\[
\lambda^2 - \left( r_0 - \delta_1 - 2\gamma \alpha^* - \frac{\alpha v^*}{(1 + \beta u^*)^2} \right) \lambda + \frac{a\alpha^* \lambda}{1 + \beta u^*} = 0,
\]

or

\[
\lambda^2 - K_{11} \lambda + K_{22} = 0,
\]

where \( K_{11} = r_0 - \delta_1 - 2\gamma \alpha^* - (\alpha v^*/(1 + \beta u^*))^2 \) and \( K_{22} = (a\alpha^*/(1 + \beta u^*))^2 \).

Due to well-known Routh-Hurwitz criterion, the characteristic (17) will have negative real roots if \( K_{11} < 0 \) and \( K_{22} > 0 \). Hence, \( r_0 - \delta_1 - 2\gamma \alpha^* - (\alpha v^*/(1 + \beta u^*))^2 < 0 \) and \( (a\alpha^* \lambda)/(1 + \beta u^*^2) > 0 \) implies \( \lambda < 0 \) since \( \delta_1, (1 + \beta u^*)^2 > 0 \). It is obvious that \( \lambda < 0 \) because \( E^* (u^*, v^*) \) will be positive interior singular point of (4). Therefore, the positive interior steady state \( E^* (u^*, v^*) \) of (4) will be asymptotically stable if \( r_0 - \delta_1 - 2\gamma \alpha^* - (\alpha v^*/(1 + \beta u^*))^2 < 0 \) and otherwise unstable.

Figure 2 shows the stability region of different singular points of (4) in \( r_0 - \beta \) parameter-space. The growth rate of prey species \( r_0 \) varies from 0 to 0.1, and the half saturation constant \( \beta \) varies from 0 to 1; the other parameters are fixed as \( \delta_1 = 0.0125, \gamma = 0.01, a = 0.8, \theta = 0.4, \) and \( \delta_2 = 0.05 \). For a very small value of \( r_0 \), that is, when \( r_0 < \delta_1 \), the trivial equilibrium \( E_0 \) is locally asymptotically stable. The stability region of \( E_0 \) is shown by blue-shaded region in the \( r_0 - \beta \) plane. In this blue-shaded region, the boundary equilibrium \( E_1 \) and the interior equilibrium \( E^* \) are unstable. The boundary equilibrium \( E_1 \) is locally asymptotically stable in the green-shaded region. But the trivial equilibrium \( E_0 \) and the interior equilibrium \( E^* \) are unstable in this green-shaded region. The red-shaded region of Figure 2 shows the locally asymptotically stable behavior of the interior equilibrium \( E^* \); that is, both the prey and predator populations can exist together in the red-shaded region. For higher value of \( r_0 \) (such as \( r_0 > 0.06 \)) and for higher value of \( \beta \) (such as \( \beta > 0.5 \)), the coexisting equilibrium \( E^* \) loses stability. The trivial equilibrium \( E_0 \) and the boundary equilibrium \( E_1 \) are unstable in the red-shaded region. All the equilibria lose stability in the white-shaded region.
responses. By considering Beddington–DeAngelis response function, the system (2) becomes
\[
\begin{align*}
\frac{du}{dt} &= r_0 u - \delta_1 u - y u^2 - \frac{a u v}{1 + \beta u + \xi v}, \\
\frac{dv}{dt} &= \frac{\alpha u v}{1 + \beta u + \xi v} - \delta_2 v.
\end{align*}
\]
(19)

The predator-prey system (19) with Beddington–DeAngelis functional response has three ecologically relevant singular points, namely, trivial singular point \( E^0 (0,0) \), boundary singular point \( E^2 (\bar{u},0) \), and an interior steady state \( E^* (\bar{u}^*, \bar{v}^*) \). Here, \( \bar{u} = (r_0 - \delta_1 / \gamma) \), and \( \bar{u}^* \) is the positive root for the following second-degree equation:
\[
\theta \gamma \bar{u}^2 - (r_0 - \delta_1) \theta - (\alpha \theta - \beta \delta_2) |\bar{u}^*-\delta_2| = 0,
\]
(20)

We have made a relative investigation of the predator-prey model (2), by considering three different functional responses, which are type-I (model (19) with \( \beta = 0, \xi = 0 \), Holling type-II (already used in model (4) with \( \xi = 0 \), and Beddington–DeAngelis. We are not showing detailed analysis of the model (2) for the different functional responses; rather, we have presented the stability scenarios for these three types of response function in Table 1. The parameters \( r_0, \theta, \) and \( \gamma \) are varied to observe the stability behavior of the system. The detailed numerical study and graphical illustrations have been shown in numerical Section 5.

3.5. Global Stability Analysis around \( E^* (u^*, v^*) \). In this subsection, we investigate the sufficient condition for global asymptotic stability of the positive coexisting singular point \( E^* (u^*, v^*) \) by establishing appropriate Lyapunov function.

**Theorem 6.** The positive interior steady state \( E^* (u^*, v^*) \) of the system (4) is globally asymptotically stable if \( (\gamma / \alpha \beta) > v^* \).

**Proof.** The interior steady state \( E^* (u^*, v^*) \) of the prey-predator model (4) fulfills the following conditions:
\[
r_0 - \delta_1 = \gamma u^* + \frac{\alpha v^*}{1 + \beta u^*}, \\
\delta_2 = \frac{\alpha u^*}{1 + \beta u^*}.
\]
(21)

With the help of these above two equations, the model (4) leads to
\[
\begin{align*}
\frac{du}{dt} &= u \left(-\gamma (u-u^*) + \frac{av^*}{1 + \beta u^*} - \frac{av}{1 + \beta u} \right), \\
\frac{dv}{dt} &= v \left(\frac{\alpha u}{1 + \beta u} - \frac{\alpha u^*}{1 + \beta u^*} \right).
\end{align*}
\]
(22)

We define the Lyapunov functional \( V(u,v) : \mathbb{R}^2_+ \rightarrow \mathbb{R} \) such that
\[
V(u,v) = V_1(u,v) + V_2(u,v),
\]
(23)

where \( V_1(u,v) = u - u^* - u^* \ln(u/u^*) \) and \( V_2(u,v) = K(\nu - \nu^* - \nu^* \ln(\nu/\nu^*)) \). Here, \( K \) is a positive constant and is defined as \( K = (1/\theta) (1 + \beta u^*) \). This particular type of Lyapunov function to study the stability of the interior equilibrium has been used widely [29, 30, 39, 40]. The Lyapunov function \( V(u,v) \) is continuous on \( \text{Int}(\mathbb{R}^2_+) \) and vanishes at \( E^* (u^*, v^*) \). Also, \( V(u,v) \) is positive definite in
In our study, the growth rate $r_0$ of prey population is one of the most important parameters. In this section, we investigate how predator-prey model behavior alters with respect to $r_0$ by using the analysis of Hopf bifurcation. We start with a theorem to study the analysis of Hopf bifurcation.

\[ \frac{dV_i}{dt} = \frac{dV_1}{dt} + \frac{dV_2}{dt}, \]

Now, differentiating (23) with respect to time $t$ and using the values of $\frac{dV_i}{dt}$ and $\frac{dV_2}{dt}$, we obtain

\[ \frac{dV}{dt} = \frac{dV_1}{dt} + \frac{dV_2}{dt}, \]

Using the relation $u^* v - \nu v^* = u^* (v - v^*) = u^* v - v^* (u - u^*)$, we obtain

\[ \frac{dV}{dt} = -\gamma (u - u^*)^2 + \frac{a \nu u^* (u - u^*)^2}{(1 + \beta u) (1 + \beta u^*)}, \]

Putting the value of $K$, we obtain

\[ \frac{dV}{dt} = -\gamma (u - u^*)^2 + \frac{a \nu u^* (u - u^*)^2}{(1 + \beta u) (1 + \beta u^*)} \]

Hence, $(\frac{dV}{dt}) < 0$ along the trajectories in $\mathbb{R}_2^+$ excluding at $E^*$ $(u^*, v^*)$ if $\gamma - a \nu u^* > 0$, that is, if $(\gamma/a \beta) > v^*$. Thus, if $(\gamma/a \beta) > v^*$, then the coexisting singular point $E^*$ $(u^*, v^*)$ is globally asymptotically stable.

4. Direction and Stability of Hopf Bifurcation

To study the stability of bifurcating limit cycle appearing through Hopf bifurcation, we use the theorem of center manifold. Center manifold theorem is an essential tool to reduce the dimension of a differential equation system in a neighborhood of a coexisting steady state (see the book by Guckenheimer and Holmes, 1983, [41]). For the predator-prey system (4), we obtain the variational matrix $J^*$ has a pair of complex conjugate eigenvalues for Hopf bifurcation.
Thus, we can investigate the model on a 2-dimensional center manifold [43].

To determine center manifold and investigate the flow here, first interpret the origin for the coordinate system to an interior singular point \( E^* (u^*, v^*) \) by expressing \( \bar{u} = u - u^* \) and \( \bar{v} = v - v^* \). For simplicity \( \bar{u} \) and \( \bar{v} \) denoted by \( u \) and \( v \), respectively, the system (4) leads to

\[
\begin{aligned}
\frac{du}{dt} &= (r_0 - \delta_1)(u + u^*) - \gamma(u + u^*)^2 - \frac{\alpha(u + u^*)(v + v^*)}{1 + \beta(u + u^*)}, \\
\frac{dv}{dt} &= \frac{\alpha \bar{v}(u + u^*)(v + v^*)}{1 + \beta(u + u^*)} - \delta_2(v + v^*).
\end{aligned}
\]

(30)

The system of (30) can be rewritten as

\[
\begin{bmatrix}
u_t \\
v_t 
\end{bmatrix} = J^*(r_0)\begin{bmatrix}u \\
v
\end{bmatrix} + \begin{bmatrix}f_1(u, v, r_0) \\
f_2(u, v, r_0)
\end{bmatrix},
\]

(31)

where \( J^*(r_0) \) defined in (16), and

\[
\begin{aligned}
f_1(u, v, r_0) &= a_1u^2 + a_2uv + a_3v^2 + \ldots, \\
f_2(u, v, r_0) &= -(a_1 + \gamma)\theta u^2 - a_2\theta uv + a_3v^2 + \ldots,
\end{aligned}
\]

(32)

where \( a_1 = -\gamma + (\alpha \beta v^*/(1 + \beta u^*)^2) \), \( a_2 = -(\alpha/(1 + \beta u^*)^2) \), and \( a_3 = 0 \).

Let the characteristic roots of the Jacobian matrix \( J^*(r_0) \) are of the form \( \lambda_{1,2} = \phi(r_0) \pm i\psi(r_0) \), where \( \phi(r_0) = (1/2)\text{tr}(J^*) \) and \( \psi(r_0) = \sqrt{\det(J^*) - \phi^2(r_0)} \). Here, the characteristics roots \( \lambda_{1,2} \) are complex conjugate if \( \det(J^*) - \phi^2(r_0) > 0 \) and the characteristic roots \( \lambda_{1,2} \) are purely imaginary; that is, \( \lambda_{1,2} = \pm i\psi(r_0) \) when \( \phi(r_0) = 0 \); that is, when \( r_0 = r_0^* \). Now set the following matrix to get the normal form of the predator-prey system (4):

\[
D = \begin{bmatrix} 1 & 0 \\
-1\hat{Q} & \hat{P}
\end{bmatrix},
\]

(33)

where \( \begin{bmatrix} 1 & -\hat{Q} \end{bmatrix} \) is the eigenvector corresponding to \( \lambda = \phi(r_0) + i\psi(r_0) \) with \( \hat{P} = (1 + \beta u^*/au^*)[(\alpha \beta u^*/v^*/(1 + \beta u^*)^2) - \gamma u^* - \phi] \) and \( \hat{Q} = ((1 + \beta u^*)\psi/au^*) \). Then, we have

\[
D^{-1} = \begin{bmatrix} 1 & 0 \\
\frac{-1}{\hat{P}} & \frac{1}{\hat{Q}}
\end{bmatrix}.
\]

By using the transformation

\[
\begin{bmatrix}u \\
v
\end{bmatrix} = D\begin{bmatrix}x \\
y
\end{bmatrix},
\]

(35)

System (30) becomes

\[
\begin{bmatrix}x_t \\
y_t
\end{bmatrix} = L_0(r_0)\begin{bmatrix}x \\
y
\end{bmatrix} + \begin{bmatrix}G(x, y, r_0) \\
H(x, y, r_0)
\end{bmatrix},
\]

(36)

where

\[
L_0(r_0) = \begin{bmatrix} \phi & -\psi \\
\psi & \phi
\end{bmatrix},
\]

(37)

with

\[
G(x, y, r_0) = \phi - \beta y u^* u^2 x^2 - \frac{\beta y u^* u^*}{(1 + \beta u^*)^2} x y + 0 \times x^2 - \frac{\beta y u^* + \phi}{(1 + \beta u^*)^2} x^2 y + 0 \times x y^2 + 0 \times y^3 + \ldots,
\]

\[
H(x, y, r_0) = -\frac{P}{Q}G(x, y, r_0) + \frac{1}{Q}\alpha g_2(x, y, r_0),
\]

\[
g_2(x, y, r_0) = \alpha \theta \left[ \frac{P}{(1 + \beta u^*)^2} - \frac{\beta v^*}{(1 + \beta u^*)^4} \right] x^2 + 0 \times x^2 + \frac{\alpha \beta Q}{(1 + \beta u^*)^2} x y
\]

\[
- \alpha \beta \theta \left[ \frac{P}{(1 + \beta u^*)^2} - \frac{\beta v^*}{(1 + \beta u^*)^4} \right] x^3 - \frac{\alpha \beta Q}{(1 + \beta u^*)^3} x^2 y + 0 \times x y^2 + 0 \times y^3 + \ldots.
\]

(38)
Now, rewriting the expression (36) in polar coordinate \((r, \zeta)\) and then expanding in the Taylor series expansion at \(r_0 = r_0^*\), we obtain

\[
\begin{align*}
\zeta_t &= \psi(r_0) + c(r_0)r^2 + \ldots = \psi(r_0^*) + \psi'(r_0)(r_0 - r_0^*) + c(r_0^*)r^2 + \ldots, \\
r_t &= \phi(r_0)r + a(r_0)r^3 + \ldots = \phi(r_0^*)(r_0 - r_0^*)r + a(r_0^*)r^3 + \ldots.
\end{align*}
\]

(39)

To investigate the stability of Hopf bifurcating periodic solution, we must have to compute the sign of the coefficient of \(a(r_0)\), where \(a(r_0)\) is given by

\[
a(r_0^*) = \frac{1}{16} \left[ G_{xxx} + G_{xyy} + H_{xx} + H_{yy} \right]_{(0,0,r_0^*)} \\
+ \frac{1}{16} \psi'(r_0^*) \left[ G_{xx}G_{yy} - H_{xy}(H_{xx} + H_{yy}) \right]_{(0,0,r_0^*)} \\
- G_{xx}H_{xx} + G_{yy}H_{yy} \right]_{(0,0,r_0^*)}.
\]

(40)

The above expression of \(a(r_0^*)\) is obtained by well-known Hopf bifurcation formula to transform (36) into Jordan form [41]. Here, the subscript denotes partial derivative, and we have

\[
G_{xxx}(0,0,r_0^*) = \frac{3\beta(\gamma u^* + \phi)}{(1 + \beta u^*)^2},
\]

\[
G_{xyy}(0,0,r_0^*) = H_{xy}(0,0,r_0^*) = G_{yy}(0,0,r_0^*) = H_{yy}(0,0,r_0^*) = 0,
\]

\[
H_{xxx}(0,0,r_0^*) = \frac{2P\beta\phi}{Qu^*(1 + \beta u^*)^2} - \frac{2a\beta\theta}{(1 + \beta u^*)^2},
\]

\[
G_{xy}(0,0,r_0^*) = -\frac{\phi}{(1 + \beta u^*)u},
\]

\[
G_{xx}(0,0,r_0^*) = \frac{2(\phi - \beta\gamma u^*)}{(1 + \beta u^*)u},
\]

\[
H_{xy}(0,0,r_0^*) = \frac{P\phi}{Q(1 + \beta u^*)u} + \frac{a\theta}{(1 + \beta u^*)^2},
\]

\[
H_{xx}(0,0,r_0^*) = \frac{2P(\phi - \beta\gamma u^*)}{Qu^*(1 + \beta u^*)} + \frac{2a\theta}{Q} \left[ \frac{P}{(1 + \beta u^*)^2} - \frac{\beta\nu^*}{(1 + \beta u^*)^3} \right].
\]

Here, \(\phi(r_0) = (1/2)tr(J^*) = (1/2)[(a\beta\nu^*/(1 + \beta u^*))^2 - \gamma u^*]\) and \((d\phi/dr_0) = (\beta\nu^*/2(1 + \beta u^*)) > 0\).

We obtain that \(\Lambda = -(a(r_0)\phi'(r_0))\).

From the above computation of \(a(r_0^*)\), the following theorem can be stated

**Theorem 8**

(i) The direction of Hopf bifurcation is supercritical, and bifurcating periodic solutions are stable if \(a(r_0^*) < 0\).

(ii) The direction of Hopf bifurcation is subcritical, and bifurcating periodic solutions are unstable if \(a(r_0^*) > 0\).

**5. Numerical Simulations**

In this section, we compute some numerical simulations regarding stability and bifurcation for the predator-prey model (4) in the vicinity of the coexisting singular point \(E^*(u^*, \nu^*)\). It is quite difficult to verify the mathematical model simulations with realistic parameter values. We have taken a hypothetical set of parameter values to illustrate our analytical findings. The set of model parameters is as follows: \(r_0 = 0.01, \delta_1 = 0.02, \gamma = 0.01, \alpha = 0.9, \beta = 0.5, \theta = 0.4, \) and \(\delta_2 = 0.05\). For this parameter, set the stability condition \(r_0 < \delta_1\) for the trivial singular point \(E_0\) is satisfied, and hence, \(E_0(0,0)\) is locally asymptotically stable. From the time series solution shown in Figure 3(a), it can be noticed that both prey and predator populations \(u(t)\) and \(v(t)\) initiate from the initial values \((0.5,0.3)\) and goes to extinct; hence, the trivial steady state \(E_0\) is stable asymptotically. Again, we have chosen the parameter values as \(r_0 = 0.03, \delta_1 = 0.01, \gamma = 0.15, \alpha = 0.9, \beta = 0.5, \theta = 0.4, \) and \(\delta_2 = 0.05\). The condition of locally asymptotic stability \(\delta_1 < r_0 < \delta_1 + (\gamma \delta_2/(\alpha \delta_3))\) of the boundary equilibrium point \(E_1(\pi,0)\) is satisfied, and hence, \(E_1(\pi,0)\) is locally asymptotically stable. Time series solution demonstrated in Figure 3(b) represents that both predator and prey populations \(u(t)\) and \(v(t)\) initiate from an initial value \((0.5,0.3)\) and go to their boundary equilibrium state, and hence, the boundary steady state \(E_1(0.1333,0)\) is stable asymptotically. Also, we have chosen the parameter set as \(r_0 = 0.03, \delta_1 = 0.02, \gamma = 0.01, \alpha = 0.9, \beta = 0.5, \theta = 0.4, \) and \(\delta_2 = 0.05\). The condition of locally asymptotic stability \(r_0 < \delta_1 - 2\gamma u^*/(\alpha v^*/(1 + \beta u^*))^2 = -0.0009 < 0\) of the coexisting singular point \(E^*(u^*, \nu^*)\) is satisfied, and hence, \(E^*(u^*, \nu^*)\) is locally asymptotically stable. From the time series analysis shown in Figure 3(c), it can be demonstrated that both prey and predators \(u(t)\) and \(v(t)\) initiate from the
The dynamical behavior around the coexisting steady state $E^*$ of the system (4) has been shown in the phase diagram (Figure 4). The system (4) is integrated by using fourth-order Runge–Kutta scheme with the parameter set $\delta_1 = 0.01$, $\gamma = 0.01$, $\alpha = 0.9$, $\beta = 0.5$, $\delta_3 = 0.05$, and $\theta = 0.4$ and varies the prey birth rate $r_0$. For $r_0 = 0.018$, the positive interior singular point is given by $E^*(0.1493, 0.0078)$. To verify Theorem 5, we find $K_{11} = r_0 - \delta_1 - 2\gamma\nu^* - (\alpha v^*/(1 + \beta u^*))^2 = -0.1113 < 0$ and $K_{21} = 0.0003 > 0$, which ensure that the asymptotic stability of coexisting state $E^*(u^*, v^*)$. Phase portrait diagram (see Figure 4(a)) demonstrates that the predator–prey model (4) is stable asymptotically at coexisting equilibrium state $E^*$ as we have chosen an initial point $E(0.5, 0.3)$ and both the species move around interior steady state $E^*(0.1493, 0.0078)$ and ultimately go to $E^*$. The system (4) enters into the limit cycle behavior from stable one, for the value of prey birth rate $r_0$ gradually increases, which means that, for increasing the size of prey population, the predator has taken more time to handle the prey species. Thus, the magnitude of the prey oscillation increases due to increasing size of the prey species. Figure 4(b) demonstrates the limit cycle oscillations around coexisting state $E^*(0.1493, 0.0400)$ for $r_0 = 0.045$, and the other parameter values are the same as in Figure 4(a). In this case, the values of the coefficients of the characteristic (17) are $K_{11} = 0.000834 > 0$ and $K_{22} = 0.001559 > 0$, which does not satisfy the conditions for locally asymptotically stability of $E^*$ in Theorem 5. As the values of $K_{11}$ and $K_{22}$ are both positive for $r_0 = 0.045$, the solutions of the system (4) are unstable, and this ensures the limit cycle oscillation around the coexisting equilibrium $E^*(0.1493, 0.0400)$.

Analytically, we have shown that the coexisting state $E^*(u^*, v^*)$ for the system (4) experiences Hopf bifurcation with respect to the prey birth rate $r_0$. By numerical simulations, Figure 5 exhibits that the predator–prey system (4) experiences Hopf bifurcation around the threshold values $r_0^* = 0.0330$. For the set of parameter values specified in this section, we obtain $\alpha d/d\alpha (\operatorname{Re}(\lambda (r_0)))|_{r_0=r_0^*} = 0.0694 > 0$. This indicates that the transcursivity condition for Hopf bifurcation is verified. Hence, the coexistence steady state $E^*(u^*, v^*)$ of the predator–prey system (4) is asymptotically stable for $r_0 < r_0^*$ and unstable for $r_0 > r_0^*$. It can also be noticed that the system (4) exhibits oscillatory behavior from the stable one, when the value of prey birth rate $r_0$ is increased gradually, and the high amplitude oscillation may lead to the crash of the species [44, 45]. Hence, the prey growth rate ($r_0$) has a critical role for the stability of the system (4).

In our study, we have not shown analytically the bifurcation for the predator–prey system (4) with respect to the conversion coefficient $\theta$ and the intraspecific coefficient $\gamma$; here, we have numerically plotted the bifurcation figure for (4) with reference to the parameters $\theta$ and $\gamma$. The bifurcation figure for (4) with respect to $\theta$ is shown in Figure 6. Here, $r_0$ is predetermined at 0.07, and the rest of parameters are the same as in Figure 5. The system (4) shows the limit cycle oscillation from equilibrium state as the bifurcation parameter $\theta$ increases. The system (4) experiences Hopf bifurcation around $\theta = 0.055$. The bifurcation diagram for the system (4) with reference to the intraspecific competition coefficient $\gamma$ is shown in Figure 7. Here, $r_0$ is predetermined at 0.07, and the rest of parameters are the same as in Figure 5. This bifurcation diagram shows that the model alters the stability from limit cycle behavior to equilibrium state as bifurcation parameter $\gamma$ increases.

Again, the dynamics around the interior singular point of the Beddington–DeAngelis-type predator–prey system (19) is shown in Figure 8. The system (19) is integrated using the Runge–Kutta scheme with parameter set $\delta_1 = 0.01$, $\gamma = 0.01$, $\alpha = 0.9$, $\beta = 0.5$, $\delta_3 = 0.05$, $\theta = 0.4$, $\xi = 0.3$, and varies the prey birth rate $r_0$. For $r_0 = 0.018$, the predator–prey system (19) has one interior positive singular point $E^*(0.1496, 0.0078)$. As shown in Table 1, we compute the value of $M_1 + M_4 = -0.0056$ and $M_1M_4 - M_2M_3 = 0.0002659$; therefore, $M_1 + M_4 < 0$ and $M_1M_4 - M_2M_3 > 0$. This ensures the stability for the interior singular point for the predator–prey system (19). The phase portrait diagram for Beddington–DeAngelis-type prey–predator system (19) shown in Figure 8(a) demonstrates that the coexisting steady state is locally asymptotically stable. The Beddington–DeAngelis-type prey–predator system (19) enters into the limit cycle behavior from stable one for increasing value
of $r_0$. Also, the phase portrait (Figure 8(b)) shows the limit cycle oscillation at an interior singular state $E^*(0.1511, 0.0405)$ for $r_0 = 0.045$.

The bifurcation diagram of the Beddington–DeAngelis-type predator-prey system (19) with respect to growth rate $r_0$ for prey species, the conversion rate $\theta$ of prey biomass to predator biomass, and the intraspecific competition rate $\gamma$ are plotted numerically in Figures 9–11, respectively. From the bifurcation diagram of the Holling Type-II predator-prey system (4) and Beddington–DeAngelis-type predator-prey system (19), it can be observed that we obtained the similar kind of dynamics of the prey and predator species except the changes in the bifurcation point. The bifurcation in Figure 9 represents that the predator-prey system (19) is asymptotically stable for lower critical values of prey birth rate $r_0$, and if $r_0$ crosses threshold value $r_0^c = 0.035$, the Beddington–DeAngelis-type predator-prey system (19) shows limit cycle behavior or oscillating behavior. The bifurcation diagram for the Beddington–DeAngelis-type predator-prey system (19) is shown in Figure 10 with respect to conversion parameter $\theta$. Figure 10 exhibits that the system alters the stability from equilibrium state to limit cycle behavior as the bifurcation parameter $\theta$ increases. The bifurcation diagram shown in Figure 11 of the Beddington–DeAngelis-type predator-prey system (19) with respect to the intraspecific competition rate $\gamma$ shows that the model alters the stability
from limit cycle behavior to equilibrium state for increasing bifurcation term \( \gamma \).

Our predator-prey system (4) with Holling type-II response function undergoes transcritical bifurcation with respect to the prey growth rate \( r_0 \), which is plotted in the bifurcation diagram in Figure 12, and the other parameter values are \( \delta_1 = 0.01, y = 0.01, \alpha = 0.9, \beta = 0.5, \) and \( \delta_2 = 0.05 \). The bifurcating parameter \( r_0 \) varied from 0 to 0.04 and is plotted \( r_0 \) along the \( x \)– axis and equilibrium prey (left panel in Figure 12) and predator (right panel in Figure 12) density along the \( y \)– axis . Black line represents the stable branch of the trivial equilibrium point \( E_0 (0, 0) \), blue-dotted line represents the stable branch of the boundary equilibrium point \( E_1 (\overline{u}, 0) \), and the green line represents the unstable branch of the boundary equilibrium point \( E_1 (\overline{u}, 0) \). The red horizontal line (left panel) and the red oblique line (right panel) indicate the stable branch of the interior equilibrium point \( E^* (u^*, v^*) \), and the magenta line indicates the unstable branch of the interior steady state \( E^* (u^*, v^*) \). There is no stable branch of the boundary steady state \( E_1 (\overline{u}, 0) \) in the plot of predator density (right panel in Figure 12) as there is no predator density in the boundary equilibrium \( E_1 (\overline{u}, 0) \). Both the prey and predator populations oscillate in the blue-shaded region.

Our predator-prey system (19) with Beddington–DeAngelis response function also undergoes transcritical bifurcation with respect to the prey growth rate \( r_0 \), which is plotted in the bifurcation diagram in Figure 13, and

Figure 6: Bifurcation diagram with respect to conversion term \( \theta \) for the predator-prey model (4) with Holling type-II response function. Rest of parameters: \( r_0 = 0.07, \delta_1 = 0.01, \gamma = 0.01, \alpha = 0.9, \beta = 0.5, \) and \( \delta_2 = 0.05 \).

Figure 7: Bifurcation diagram with respect to intraspecific competition between prey species \( \gamma \) for predator-prey system (4) with Holling type-II response function. Rest of the parameters \( r_0 = 0.07, \delta_1 = 0.01, \alpha = 0.9, \beta = 0.5, \) and \( \delta_2 = 0.05 \).
Figure 8: Phase portrait diagram of the predator-prey system (19) with Beddington–DeAngelis-type response function at the coexisting singular point $E^\ast$ (green filled circle) with initial values $[x(0), y(0)] = [0.5, 0.3]$. Black line is the predator isocline, and the red curve is the prey isocline. (a) $r_0 = 0.018$; (b) $r_0 = 0.045$, and rest of the parameters are specified as $\delta_1 = 0.01$, $\gamma = 0.01$, $\alpha = 0.9$, $\beta = 0.5$, $\theta = 0.4$, $\delta_2 = 0.05$, and $\xi = 0.30$.

Figure 9: Bifurcation diagram with respect to birth rate of prey species $r_0$ of the predator-prey system (19) with Beddington–DeAngelis-type response function. Other parameter values are specified as $\delta_1 = 0.01$, $\gamma = 0.01$, $\alpha = 0.9$, $\beta = 0.5$, $\theta = 0.4$, $\delta_2 = 0.05$, $\theta = 0.4$, and $\xi = 0.30$.

the other parameter values are $\delta_1 = 0.01$, $\gamma = 0.01$, $\alpha = 0.9$, $\beta = 0.5$, $\theta = 0.4$, $\delta_2 = 0.05$, and $\xi = 0.30$. The bifurcating parameter $r_0$ varied from 0 to 0.04 and is plotted $r_0$ along the $x$–axis and equilibrium prey (left panel in Figure 13) and predator (right panel in Figure 13) density along the $y$–axis. Black line represents the stable branch of the trivial equilibrium point $E^0(0, 0)$, blue-dotted curve represents the stable branch of the boundary equilibrium point $E^1(\tilde{u}, 0)$, and the green line represents the unstable branch of the boundary equilibrium point $E^1(\tilde{u}, 0)$. Red curve indicates the stable branch of the interior equilibrium point $E^\ast(\tilde{u}^\ast, \tilde{v}^\ast)$, and the magenta line indicates the unstable branch of the interior steady state $E^\ast(\tilde{u}^\ast, \tilde{v}^\ast)$. There is no stable branch of the boundary steady state $E^1(\tilde{u}, 0)$ in the plot of the predator density (right panel in Figure 13) as there is no predator density in the boundary equilibrium $E^1(\tilde{u}, 0)$. Both the prey and predator populations oscillate in the blue-shaded region. From the transcritical bifurcation for Holling type-II and Beddington–DeAngelis-type response function, we can conclude that density of prey species increases (see the red curve of Figures 12 and 13) for Beddington–DeAngelis-type response function. The red curve (left panel in Figure 13) is increasing as the value of $r_0$ is increasing; that is, the density of prey species is increasing as the prey growth rate $r_0$ is increasing. To better understand the increasing pattern of prey species for Beddington–DeAngelis-type response function, we put a portion of the red curve in the dash-dotted box, which has been zoomed and plotted in the inset (left panel of Figure 13).
By constructing suitable Lyapunov function, we have analytically shown that (see the Section 3.5) the coexisting steady state $E^* (u^*, v^*)$ for the predator-prey system (4) is globally asymptotically stable for $(\gamma/a\beta) > v^*$. For the following parameter values $r_0 = 0.045$, $\delta_1 = 0.01$, $\gamma = 0.15$, $\alpha = 0.9$, $\beta = 0.5$, $\theta = 0.4$, and $\delta_2 = 0.05$, we obtain the interior equilibrium is $E^* (0.1493, 0.0151)$. Here, $(\gamma/a\beta) = 0.3333$ and $v^* = 0.0151$, so it can be verified that $(\gamma/a\beta) > v^*$ for the assumed set of parameter value. Therefore, the interior steady state $E^* (0.1493, 0.0151)$ is globally asymptotically stable. Now, we have numerically integrated the predator-prey system (4) using the Runge–Kutta scheme with the assumed parameter values and considered different initial values $(0.4, 0.1)$, $(0.2, 0.2)$, $(0.4, 0.1)$, and $(0.1, 0.3)$. The phase diagram of this solution is plotted in Figure 14. All the solution trajectories for different initial values converge to an interior singular point $E^* (0.1493, 0.0151)$. This numerically ensures that the interior steady state $E^*$ for the model (4) is globally asymptotically stable.

6. Discussion

Mathematical modeling is an important powerful tool to better understand the interplays between various species that allow us to continue biodiversity and ecosystem in essence.
In literature, a series of mathematical models has been explored to understand the dynamics for prey-predator model with different kinds of functional responses [12]. In this paper, we investigate a very simple mathematical model with different kinds of functional responses including Beddington–DeAngelis-type response function. Other parameter values are specified as $\delta_1 = 0.01$, $\gamma = 0.01$, $\alpha = 0.9$, $\beta = 0.5$, $\theta = 0.4$, and $\delta_2 = 0.05$.

The solutions of the appraised predator-prey model are showed to be bounded, which indicate none of the species would go at its infinite. To study the kinetics of (4), we discussed the existence of biologically meaningful singular points and their stability including global asymptotic stability. The singular points are identified using prey and predator isocline and found three biologically meaningful steady states of the system (4). The boundary steady state $E_i(\bar{u}, \bar{v})$ shows oscillatory dynamics to the present system (4) as the stability for the model around $E_i$ indicates the infeasibility for the trivial equilibrium point $E_0$. The boundary steady state $E_i(\bar{u}, \bar{v})$ is feasible if $r_0 > \delta_1$; that is, growth rate of the prey species is higher than the decay rate.
librium point of predator population. Moreover, we notice that the population by predator population lowers the density of equilibrium. The results demonstrate that birth suppression of prey population plays a critical role to destabilize the predator-prey system (4). The theoretical and numerical illustrations for our present paper describe that a simple predator-prey system can have rich dynamics which can be observed in the ecological system. Numerical simulations are employed for the complexity of theoretical findings, where achievable. For the aptitude suggestions of our investigation, there is a permanent need for observational work to test these forecasts. It can also be observed that the illustrations represented in our manuscript should be examined from qualitative, preferably, than quantitative view point. Our study is based on hypothetical set of parameter values and not based on any particular case as we have no experimental data in our hand, but this investigation might be useful for ecologists who are performing their observation in similar field on the basis of experimental data.

**Data Availability**
All relevant data are included within the paper.

**Conflicts of Interest**
The authors declare that they have no conflicts of interest.

**Acknowledgments**
The work of J. J. Nieto has been partially supported by the Agencia Estatal de Investigación (AEI) of Spain, cofinanced by the European Fund for Regional Development (FEDER) corresponding to the 2014–2020 multiyear financial framework (project MTM2016-75140-P) and Xunta de Galicia (grant ED431C 2019-02). The work of Subhas Khajanchi was supported by the Science and Engineering...
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