GRB980923. A BURST WITH A SHORT DURATION HIGH-ENERGY COMPONENT

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ABSTRACT

The prompt emission of gamma-ray bursts (GRBs) is usually well described by the Band function: two power laws (PLs) joined smoothly at a given break energy. In addition to the Band component, a few bursts (GRB941017, GRB090510, GRB090902B, and GRB090926A) show clear evidence of a distinct high-energy spectral component, which in some cases evolves independently from the prompt keV component and is well described by a PL, sometimes with a cutoff energy; this component is found to have long duration, even longer than the burst itself for all four bursts. Here we report the observation of an anomalous short duration high-energy component in GRB980923. GRB980923 is one of the brightest GRBs observed by BATSE. Its light curve is characterized by a rapid variability phase lasting ~40 s, followed by a smooth emission tail lasting ~400 s. A detailed joint analysis of BATSE (Large Area Detector and Spectroscopy Detector) and EGRET TASC data of GRB980923 reveals the presence of an anomalous keV to MeV component in the spectrum that evolves independently from the prompt keV one. This component is well described by a PL with a spectral index of −1.44 and lasts only ~2 s; it represents one of the three clearly separated spectral components identified in GRB980923, the other two being the keV prompt emission, well described by the Band function and the tail, well fit by a Smoothly Broken Power Law.

Key words: acceleration of particles – astroparticle physics – gamma-ray burst: general – gamma-ray burst: individual: GRB980923 – radiation mechanisms: non-thermal

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1. INTRODUCTION

Gamma-ray bursts (GRBs) are the most energetic sources in the universe, with a total isotropic-equivalent radiated energy in the range \(10^{48}–10^{55}\) erg. They are characterized by a brief and intense flash, the so-called prompt emission, observed in γ- and X-ray energy bands, followed by a long-lived afterglow emission, radiated in the X-ray band and below. The GRB prompt emission spectra are typically best fit by the “Band function”: two power laws (PLs) joined smoothly at a given break energy (Band et al. 1993), whose low-energy and high-energy photon indices, α and β, have median values of −1 and −2.3, respectively (Preece 2000; Kaneko et al. 2006). In addition to the Band component, a few bursts show clear evidence of a distinct high-energy (HE) spectral component. The first GRB showing such a characteristic was GRB941017, for which an extra multi-MeV spectral component has been observed (González et al. 2003). This component lasts longer (200 s) than the keV component (T90 = 77 s), from which it evolves independently in time; it is well fit by a PL with spectral index −1.0 and carries two-thirds of the energy fluence.

With the launch of the Fermi gamma-ray space telescope (Atwood et al. 2009; Meegan et al. 2009) three additional bursts with a clear distinct HE spectral component besides the Band function have been observed: GRBs 090910 (Ackermann et al. 2010), 090902B (Abdo et al. 2009), and 090926A (Ackermann et al. 2011). This component is usually well fit by a hard PL that dominates at higher energies (Granot 2010) and has a long duration, longer than the duration of the burst itself.

Specifically, GRB090910, characterized by T90 = 2.1 s, presents an extra PL component with a spectral index −1.62 during the prompt phase; this short duration HE emission was delayed by ~0.1 s with respect to the onset of the Band component and was followed by an extended HE emission, lasting ~200 s and well fit by a PL (Ackermann et al. 2010; de Pasquale et al. 2010).

GRB090902B, characterized by \(T_{90} = 21.9\) s, shows a hard additional component during the prompt phase which is well fit by a single PL with spectral index −2.1 and lasts ~1000 s. Furthermore, a spectral feature at energies < 50 keV is evident; this is consistent with an extrapolation of the HE PL emission down to those energies (Abdo et al. 2009).

The last Fermi burst, GRB090926A, shows an extra HE component that is very significant at the time of a narrow pulse (lasting less than 1 s) observed in the prompt emission light curve; the HE emission is also present at later times, with an overall duration of hundreds of seconds (see also Section 4.1). The fit of the narrow pulse spectrum requires, as seen for the other bursts, an extra hard PL component besides the Band one, but for the first time a spectral break has been detected in this PL around 1.41 GeV. After this pulse the HE emission is well described by a simple PL (Ackermann et al. 2010).

There are two main classes of models that have been proposed to describe the HE γ-ray emission: leptonic and hadronic. Between the leptonic models, the most investigated scenarios consider inverse Compton (IC) and synchrotron self-Compton (SSC) emission processes in different locations of the relativistic jet from which the GRB is generated. Two regimes for accelerating electrons are considered: They can be accelerated in external (as the jet interacts with the circumburst medium) and internal (within the jet as the Lorentz factor of the flow varies) shocks. Different seed photon populations for the IC scattering have been discussed in detail in GRB internal shocks (Papathanassiou & Mészáros 1996; Pilla & Loeb 1998; Panaitescu & Mészáros 2000), forward shocks (Sari et al. 1998; Totani 1998a; Waxman 1997; Panaitescu & Mészáros 1998;...
GRB090902B, however, the keV–GeV spectrum detected in (2004). They considered a beam of ultra HE neutrons produced HE emission observed in GRB941017 by Dermer & Atoyan that has a long duration in comparison with the burst producing the anomalous HE component. Internal and external radiation photons; this neutron beam can in the photopion interactions of ultra-relativistic protons with producing the anomalous HE component such as the one of GRB941017. Also, synchrotron emission propagating into an ambient with a sufficiently low density has been proposed to explain the HE emission observed in GRB090510 (He et al. 2011) and GRB090902B (Liu & Wang 2011). More recently, Fraija et al. (2012) and Veres & Mészáros (2012) explored SSC processes from forward and reverse shocks as possible mechanisms to produce the HE emission and concluded that, depending on the equipartition parameters for the magnetic field and the electron energy in the forward and reverse shock, one or the other could develop to explain the HE component.

Within the hadronic models, the two most investigated scenarios consider the following, respectively: (1) synchrotron radiation from ultra HE protons accelerated in the relativistic jet (Vietri 1997; Totani 1998a, 1998b; Razzano et al. 2009, 2010); (2) synchrotron and IC emission from secondary electron–positron pair cascades triggered by photopion interactions of shock accelerated protons with low-energy photons coming from the prompt emission or from an external radiation field (Waxman & Bahcall 1997; Böttcher & Dermer 1998). Asano et al. (2009) investigated both scenarios to explain the HE emission observed in GRB090510 and found that, in both cases, to produce the HE component of this burst the proton injection isotropic-equivalent luminosity must be larger than $10^{55} \text{erg s}^{-1}$, thus larger than the $\gamma$-ray luminosity; in the case of GRB090902B, however, the keV–GeV spectrum detected in GRB 090902B is well explained with a comparable energy in protons and $\gamma$ rays (Asano et al. 2011).

An additional process has been proposed to explain the HE emission observed in GRB941017 by Dermer & Atiyon (2004). They considered a beam of ultra HE neutrons produced in the photopion interactions of ultra-relativistic protons with internal and external radiation photons; this neutron beam can undergo further interactions with external photons to produce hyper-relativistic electrons that radiate energy by synchrotron, producing the anomalous HE component.

Most of the models described consider an extra HE component that has a long duration in comparison with the burst $T_90$ and that is well described by a hard PL, with a photon index almost constant in time.

Within this context an interesting case of study is GRB980923, one of the brightest bursts detected by the Burst And Transient Source Experiment (BATSE), characterized by an anomalous short duration HE emission during the prompt phase. Some evidence for the presence of this extra HE component was first reported by González (2005), who analyzed joint data from the BATSE Large Area Detector (LAD) and the Energetic Gamma-Ray Experiment Telescope (EGRET) calorimeter, the Total Absorption Shower Counter (TASC), and found that an extra PL, besides the Band function, is needed to correctly reproduce the spectrum of this burst from 13 s to 46 s. Here we extend the spectral analysis of GRB980923 using, besides the data from the BATSE-LAD and EGRET-TASC, data from the BATSE Spectroscopy Detector (SD), gaining time resolution and energy coverage (see Section 3). This allows us to determine with more accuracy the time duration as well as the spectral index of the HE component. In particular, time-resolved spectral analysis revealed that this HE component only lasts $\sim 2$ s, in contrast with the overall long duration of the extra HE component of the other GRBs. This component represents one of the three clearly separated spectral components identified in GRB980923, the other two being the prompt emission, well described by the Band function and a tail, lasting 400 s and well fit by a PL (Giblin et al. 1999).

The paper is organized as follows: In Section 2 we summarize the main observational properties of GRB980923. In Section 3 we briefly describe BATSE and EGRET-TASC and the set of data chosen; we also explain how the analysis of these data was performed. In Section 4 we present the results. In Section 5 we present our conclusions.

2. GRB980923

GRB980923 was observed by BATSE at 20:10:52 for 33.02 s (as determined by BATSE T90) from the galactic coordinates $l = 293:08$ and $b = -30:89$. The burst direction was 124.4 with respect to the pointing-axis direction of the Compton Gamma Ray Observatory (CGRO), placing the burst outside the $\sim 1$ sr field of view of the Imaging Compton Telescope (COMPTEL) and the EGRET spark chamber. With a fluence of $4.5 \times 10^{-5} \text{erg cm}^{-2}$, GRB980923 is the burst with the third highest fluence in the BATSE catalog (Preece 2000).

The prompt emission light curve of GRB980923 was observed by BATSE-LAD and BATSE-SD has two main episodes (see Figure 1, three top panels). The first episode lasts about $\sim 14$ s, then the count rate drops drastically, for about 2 s, to start the second episode, whose duration is $\sim 20$ s. A single Band function is not sufficient to describe the whole prompt emission from 20 keV to 200 MeV, as shown in the joint analysis of data from BATSE-LAD and EGRET-TASC performed by González (2005), Kaneko et al. (2008), and González et al. (2009).
In particular, González (2005) found that to fit the second episode an extra PL is needed.

The prompt emission of GRB980923 is followed by a long and smooth tail, whose duration is reported to be ∼400 s (Giblin et al. 1999). This tail was specifically studied by Giblin et al. (1999) using BATSE data. Its time-integrated spectrum is well described by a Smoothly Broken Power Law (SBPL), as well as by the Band function. Giblin et al. (1999) identified a separate emission component related to the tail starting at ∼40 s when the spectral shape changes dramatically. The temporal evolution of the energy break of the SBPL showed that if the tail starts at 32.109 s after the trigger, the evolution of its spectrum corresponds to the evolution of a synchrotron cooling break in the slow-cooling regime and the transition from fast cooling could happen on timescales comparable to the duration of the burst.

Therefore, if the synchrotron emission from external shocks explains the tail then, at least for some bursts, the afterglow may begin during the prompt phase (Giblin et al. 1999).

3. BATSE AND EGRET-TASC DATA

In this work, data from BATSE and EGRET-TASC were used. A brief description of these instruments and of the set of data chosen is given below, together with some details about how the spectral analysis of the data was performed.

BATSE (for a detailed description see Fishman et al. 1989; Pendleton et al. 1995; Preece 2000 and Kaneko et al. 2006) consisted of eight separate identical detector modules located at the corners of the CGRO spacecraft on the faces of a regular octahedron. Each module had two NaI(Tl) scintillation detectors coupled with photomultiplier tubes (PMTs): a LAD optimized for sensitivity and directional response and a SD optimized for energy coverage and resolution. Each detector had different energy capabilities: the LAD had a constant energy range of 0.02–1.9 MeV, while the SD had an adjustable energy range between 0.01 and 100 MeV depending on the PMT’s gains. Even though the SD was sensitive over a broader energy range, the LAD had a collecting area 16 times bigger than the SD’s. Each detector module was independent. Data were processed and accumulated to construct various data types with different energy resolutions and accumulation times in the data processing unit. In the analyses of González et al. (2003) and González (2005), only LAD data were used because of the bigger collecting area. In this paper, we also include SD data to increase the energy range with finer time resolution than EGRET-TASC. In the case of the LAD, we use data from detector 7 (LAD7) in the Continuous (CONT) format that contains 16 energy channels with an accumulation time of 2.048 s. In the case of the SD, we use data from detectors 3 (SD3) and 7 (SD7) in the Spectroscopy High Energy Resolution Burst format, which contains spectra in 256 energy channels with an accumulation time of 512 ms. SD7 and SD3 data allow us to extend the analysis to lower (32 keV) and higher energies (27 MeV), respectively.

TASC is a monolithic 76 × 76 × 20 cm³ scintillation calorimeter formed by 36 NaI(Tl) blocks optically coupled. It was viewed by two groups of eight interleaved PMTs, each feeding a pair of Pulse Height Analyzers (PHAs). One pair of PHAs processed low-energy events (1–200 MeV), while the other processed HE events (0.02–30 GeV). Although the TASC was part of the EGRET instrument measuring the energy of each useful event triggering the spark chambers (mainly done with the HE PHA), it was also an independent detector sensitive to gamma rays and charged particles (using the low-energy PHA).

As an independent detector, the accumulation time of 32.768 s for the low-energy spectrum was the normal and continuous mode of data acquisition and was called the solar mode. The TASC data used for GRB980923 consist of only two solar spectra with 229 energy channels, starting 19.6 s before the BATSE trigger and covering almost twice the burst duration (for a detailed description, see González et al. 2009).

The light curves of the emission observed by LAD7, SD7, SD3, and TASC are shown in Figure 1. General features are seen in all the BATSE light curves: two main episodes separated by a drop in the count rate at ∼14 s with a duration of ∼2 s (see also Section 2). The drop is not observed in TASC data because of the 32.768 s accumulation time of the spectra. The finer time resolution of the SD7 and SD3 data allows us to observe the structure of both episodes. In particular, an intense peak at ∼20 s is more evident from the lowest to the highest BATSE energies. The long tail is evident in LAD data.

The spectral fitting was performed with the software developed by the BATSE team, RMFIT.4 For a detailed description of the spectral analysis, see González et al. (2009); here, only details relevant to the analysis of GRB980923 are mentioned. An important step of the analysis is the background subtraction. The time dependency of the background was given mainly by the spacecraft position with respect to Earth’s magnetic field. We modeled background counts with a polynomial of the fourth order on each energy bin using spectra for ∼500 s before and after the burst trigger. The background model was checked against data corresponding to 15 earlier/later orbits when the spacecraft was located at the same geomagnetic rigidity. Background subtracted spectra of LAD7, SD3, SD7, and TASC data were fitted using the detector’s responses. When a joint fit of BATSE and TASC data was performed, BATSE data were binned in time to match the time resolution of TASC data; when only BATSE data were used, SD data were binned to match the time resolution of LAD data. The energy ranges used in the fit for each data set are the same given for the light curves; see Figure 1. The spectral fitting was performed with the following photon flux models: Band function, and Band function plus a PL and a SBPL. As in González et al. (2009), normalization factors between data sets were introduced to account for errors in the calculated effective area of the detectors because of underestimation of the CGRO mass model. The value of the normalization factor depends mostly on the zenith angle of the burst. If the detector responses were perfectly known, the normalization factor would be one. The factors normalize TASC and SD data with respect to LAD data, and it is the same for all energies in a given data set. However, in our analysis it varies from time bin to time bin because it also accounts for the mismatching of time binning between data sets. These normalization factors had similar values for both SD detectors. The TASC normalization factor is, as expected, higher mainly because of the greater accumulation time of TASC with respect to the other detectors (see Section 4).

4 R. S. Mallozzi, R. D. Preece, & M. S. Briggs, “RMFIT: A Light curve and Spectral Analysis Tool,” © Robert D. Preece, University of Alabama in Huntsville.
Figure 2. Spectral fitting for different time intervals as defined in Section 4. LAD7 (black), SD7 (blue), SD3 (green), and TASC (red) are shown. For the purpose of the plot, but not for the spectral fit, the data are binned in energy to give at least 2σ significance over background. Solid curves show model fits to the data using the parameters given in Table 1 and the spectral model described in the text. The upper limits correspond to 2σ deviation from the background. The two spectral components, the Band function at lower energies and the higher-energy power law, are shown as dashed lines. The intervals excluding the time interval from 19.5 to 21.5 s are well adjusted only with a Band function, otherwise a PL function is required to fit the highest energies.

(A color version of this figure is available in the online journal.)

4.1. Spectral Fit and Identification of the HE Component

We fit jointly LAD7-SD3-SD7-TASC data for the time intervals 0–13 s and 13–33 s determined by BATSE. The first time interval is fitted with a Band function, while the second one also requires a HE-PL function described by $A_{PL}[E(\text{keV})/10\text{ MeV}]^\gamma$. The data and the spectral fits are shown in Figure 2 (two top panels), with the best-fit parameters and the normalization factors given in the first two columns of Table 1. The normalization factors between the LAD and SD detectors for both time intervals are very similar because the time coverage almost matches. However, the TASC normalization factors are different mainly because they account for the 18 s and 12 s of extra TASC time coverage compared to LAD time coverage in the first and second time interval, respectively. Also, the normalization factors are calculated considering the best fit to all data sets simultaneously. This and the contribution of the PL (when required) are small when compared with the corrections for the time bin mismatching, but they explain the difference with the values given by González et al. (2009).

The energy dependency, as well as the correct election of the normalization factors, is confirmed by the fact that SD and TASC data below 10 MeV smoothly continue the lower-energy component described by the Band function in both time intervals (see Figure 2).

In order to better characterize the HE component and to determine with greater accuracy its duration, we look at the light curves as a function of energy for all detectors; see Figures 3–6. The peak at $\sim20$ s is evident from 8 keV up to 28 MeV as observed by SD7 and SD3, respectively. It evolves with energy differently from the whole light curve, which dims at higher energies. This is consistent with the TASC data: The first time interval dims at higher energies, contrary to the second one, which includes the HE component. We fit jointly LAD-SD data for 2 s time intervals from 0 to 33 s and find that only the time interval corresponding to the peak from 19.5 to 21.5 s requires a PL besides a Band function to fully describe the data. The data and the spectral fits are shown in Figure 2 (lower panels), with the best-fit parameters given in Table 1; the time evolution of the Band spectral indices is shown in Figure 7. It can be seen that, except for the value of the low-energy spectral index $\alpha$ at $\sim14$ s and within the error bars, a smooth time evolution is observed in both the Band spectral indices through the whole time period from 0 to 33 s. This suggests that the prompt emission described by the Band function comes from the same radiation process; the different value of $\alpha$ at $\sim14$ s could indicate the presence of a distinct component starting during

Figure 3. LAD7 light curves in different energy bands (energy increasing from top to bottom). A peak at 20 s becomes more evident at the highest energies.
Table 1
Spectral Fitting Parameters of the Differential Photon Flux

| Time from BATSE Trigger (s) | 0 to 13 | 13 to 33 | 13 to 19.5 | 19.5 to 21.5 | 21.5 to 33 |
|-----------------------------|---------|----------|------------|-------------|------------|
| Band GRB function parameters |         |          |            |             |            |
| $A$ ($\times 10^{-2}$ photon s$^{-1}$ cm$^{-2}$ keV$^{-1}$) | 14.6$^{+0.1}_{-0.2}$ | 24.7$^{+0.2}_{-0.3}$ | 12.3$^{+0.2}_{-0.2}$ | 29.5$^{+0.1}_{-0.1}$ | 32.1$^{+0.3}_{-0.2}$ |
| $E_{\text{peak}}$ (keV) | 345.4$^*$ | 371.3$^*$ | 419.15$^*$ | 419.15$^*$ | 346.32$^*$ |
| $\alpha$ | $-0.476^{+0.01}_{-0.01}$ | $-0.573^{+0.002}_{-0.002}$ | $-0.549^{+0.021}_{-0.021}$ | $-0.554^{+0.009}_{-0.009}$ | $-0.563^{+0.014}_{-0.007}$ |
| $\beta$ | $-2.6^{+0.04}_{-0.04}$ | $-3.06^{+0.09}_{-0.11}$ | $-2.48^{+0.07}_{-0.07}$ | $-4.01^{+0.07}_{-0.03}$ | $-3.08^{+0.07}_{-0.07}$ |
| High-energy power-law parameters |          |          |            |             |            |
| $A_{\text{HE PL}}$ ($\times 10^{-5}$ photon s$^{-1}$ cm$^{-2}$ keV$^{-1}$) | ... | $0.7^{+0.2}_{-0.2}$ | ... | $5.6^{+0.8}_{-0.8}$ | ... |
| $\gamma$ | ... | $-1.28^{+0.26}_{-0.33}$ | ... | $-1.44^{+0.07}_{-0.07}$ | ... |
| Probability | ... | $6.20 \times 10^{-7}$ | ... | $8.12 \times 10^{-11}$ | ... |
| Normalization factors |         |          |            |             |            |
| SD7 | 1.28 | 1.31 | 1.21 | 1.44 | 1.36 |
| SD3 | 1.23 | 1.30 | 1.05 | 1.51 | 1.39 |
| TASC | 0.76 | 0.94 | ... | ... | ... |

Notes. The first six rows contain the best spectral fit parameters for different time intervals and including (Columns 2 and 3) or not (Columns 4–6) TASC data. The seventh row shows the probability that the improvement in $\chi^2$ from the addition of the HE PL in the fit is due to chance, as determined by the $\chi^2$ test. The last three rows are the normalization factors relative to LAD7 (see Sections 3 and 4) used in the fit.

Figure 4. SD3 light curves in different energy bands (energy increasing from top to bottom). A peak at 20 s becomes more evident at the highest energies.

Figure 5. SD7 light curves in different energy bands (energy increasing from top to bottom). A peak at 20 s becomes more evident at the lowest and highest energies.

The prompt phase (see Section 4.2), which could be noticeable because of the dimming of prompt photons at $\sim 14$ s (see the three upper panels in Figure 1). Concerning the PL component, the value of the probability that the improvement in $\chi^2$ from the addition of the HE PL in the fit is due to chance, as determined by the $\chi^2$ test considering the time interval from 19.5 to 21.5 s is four orders of magnitude (six when TASC data are included) smaller in comparison with the one obtained for the time interval between 13 s and 33 s (see Table 1). We also fit smaller time intervals included in the period from 19.5 s to 21.5 s using only SD data, and found that the smallest probability that the improvement is due to chance is found when considering the.
Figure 6. TASC light curves in different energy bands (energy increasing from top to bottom). The time interval including the peak at 20 s, as observed by BATSE, remains significant up to hundreds of MeVs.

Table 2

| GRB Name    | HE Duration (s) | PL Index | Tail |
|-------------|-----------------|----------|------|
| (1) GRB 980923 | 2              | -1.44 ± 0.07 | Yes |
| (2) GRB 941017 | ~200           | ~-1      | Yes |
| (3) GRB 090510 | ~200           | -1.62 ± 0.03 | Yes |
| (4) GRB 090902B | ~1000         | -1.93 ± 0.01 | No  |
| (5) GRB 090926A | ~4800         | -1.79 ± 0.02 | Yes |

Notes.

a References for the HE duration: (1) see text; (2) González et al. 2003; (3) Ackermann et al. 2010; (4) Abdo et al. 2009; (5) Ackermann et al. 2011.

b References for the PL spectral index: (1) see text; (2) González et al. 2003; (3) Ackermann et al. 2010; (4) Abdo et al. 2009; (5) Ackermann et al. 2011.

c References for the tail: (1) Giblin et al. 1999; (2) González 2005; (3) Guiriec et al. 2009; (5) Noda et al. 2009, Golenetskii et al. 2009.

whole 2 s time period. This is an indication of the fact that the PL was always required along the 2 s time period. Therefore, the HE component of GRB980923, extending from a few keVs to hundreds of MeVs, starts at 19.5 s and lasts only 2 s; GRB980923 represents, then, the first case of a burst showing an extra HE PL component of short duration (see Table 2 for a comparison with the other bursts). The energy flux of this HE component is $22 \times 10^{-6}$ erg s$^{-1}$ cm$^{-2}$ in the energy range from 2 to 200 MeV.

A burst that presents some similarities with GRB980923 is GRB090926A. This source also is characterized by a HE emission associated with a short spike (lasting less than 1 s) observed in the prompt emission light curve, although a HE emission is also observed after the spike (see Section 1). The HE emission associated with the narrow pulse is described by a PL, having a different spectral index with respect to the one observed at later times and extends to the lowest energies, similarly to GRB980923. Moreover, a spectral break and the extension to the lowest energies are evident only when fitting the spike spectrum. These observational properties suggest that we are in the presence of two distinct HE components: one of short duration that could have been produced by the same physical mechanism responsible for the HE emission observed in GRB980923, and one of long duration (Sacahui et al. 2012). Other common characteristics between GRB980923 and GRB090926A are a drop in the count rate observed before the spike and the presence of a tail (see Table 2).

4.2. Spectral Fit and Implications for the Tail

We do the fit of time-resolved prompt emission spectra also using a SBPL for comparison with Giblin et al. (1999). Because of the low statistics at energies above 1 MeV, SBPL seems to be preferred by Giblin et al. (1999). However, we do not find in the data any preference for using SBPL or Band. The following discussion is independent of the fitting function used.

The time evolution of the low-energy and high-energy indices for the BAND and SBPL functions ($\alpha$, $\beta$, and $\beta_{\text{SBPL}}$, respectively) is shown in Figure 7. It can be seen that the values of $\alpha_{\text{SBPL}}$ and $\alpha_{\text{SBPL}}$ at 14 $\text{s}$ do not follow the general trend. In particular, the value of $\beta_{\text{SBPL}}$ is consistent with the one of the tail as given by Giblin et al. (1999), a possible indication of the presence of the tail already at 14 $\text{s}$. Furthermore, a higher dispersion of the $\alpha$ and $\alpha_{\text{SBPL}}$ values after $\sim$14 $\text{s}$ is noticeable.
(see Figure 7), and this could be interpreted as further evidence of the presence of the tail from 14 s. In fact, the time-integrated spectrum of the tail can be described with a SBPL with break energy around 200 keV (Giblin et al. 1999); therefore, if the tail is present from 14 s, it will affect the determination of the low-energy spectral index of the Band function, which has a peak energy around 400 keV. So, it is possible that the tail starts at \( \sim 14 \) s from the burst trigger instead of 32.109 s as reported by Giblin et al. (1999; see also Section 2). In this case, its behavior is still similar to that of the afterglows at lower energies, but the new estimation of its starting time sets the cooling regime somewhere in the transition between fast and slow cooling.

Giblin pointed out that the starting time of the tail could happen before 32 s; however, the fact that the tail could start at \( \sim 14 \) s implies that there could be a causality relationship between the tail and the HE PL component.

Surprisingly, also almost all the other bursts characterized by an extra HE PL component present a tail (see Table 2). The only exception is GRB 090902B, for which no tail has been reported in the literature, although a weak emission similar to a tail is exception is GRB 090902B, for which no tail has been reported in the literature, although a weak emission similar to a tail is

| Spectral Component | Photon Flux Function | Peak Energy | Break Energy | Duration | Start Time | Model |
|--------------------|----------------------|-------------|--------------|----------|------------|-------|
| Main prompt emission | Band function | 400 keV | \( \cdots \) | 33 s | 0 | Internal shocks |
| Tail | SBPL | \( \sim 200 \) keV | 400 s | 14 s | Synchrotron external shocks |
| High-energy component | PL | \( >200 \) MeV | \( \cdots \) | 2 s | 19.5 s | Unknown |

5. CONCLUSION

We have studied the prompt keV–MeV emission of GRB980923 as observed by BATSE and TASC. We have shown that GRB980923 presents three different spectral components: the main component described by a Band function with peak energy of the order of 400 keV, a tail component described by a SBPL with a break energy of the order of 200 keV and a HE component described by a PL. Table 3 summarizes the main characteristics of each spectral component. The duration of this HE component is shorter than the duration of the burst itself, which is different compared with the other GRBs that present a HE component distinct from the Band function. Also, this HE component is observed in the middle of the burst episode, not at the end or beginning. HE emission with a much shorter duration than the prompt emission clearly challenges the actual models, in which a HE component has to last as long as the prompt phase or longer.

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