The Mathematical Model of Arc Discharge in Metal Vapors at Active Gases over Crucible for Technological Process of Electron Beam Deposition of Ceramic Coatings

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ABSTRACT

The mathematical model of arc discharge in the metal vapors, propagated in the soft vacuum in the medium of active gases, is presented in this chapter. Such type of discharge is widely used in advanced electron beam technologies for obtaining the coating of new types from nanostructurized materials, especially ceramics coating. As electron beam sources for evaporation of refractory metals in this technological process the high voltage glow discharge electron guns are widely and effectively used. But the aim of applying of additional low-voltage arc discharge under the crucible is stimulating and maintaining the chemical reactions between metal vapors and residual gas in the vacuum technological chamber. In formed model for calculation of electric field distribution the analytical solving of Poisson equation is used, and the spatial distribution of discharge current density is defined on the base of the equation of uncontineously of electrons and ions flows. All analytical relations for auxiliary geometry of electrodes system with cylindrical crucible and ring electrode with positive potential over it have been obtained. Such electrodes system is standard for electron beam installations, designed to deposition of chemically complex ceramic coatings. The simulation results shown, that for power of electron beam range of nearly 10 kW and greater, the pressure in technological chamber range of few Pa and potential on electrode, located over crucible, 70 – 100 V, the current density of arc discharge is in range 0.05 – 0.1 A/cm\(^2\). Such value of current density is generally enough to maintaining the chemical reaction between the metal vapors and active gas for obtaining the nanostructurized coatings with good steheometry. Obtained simulation results may be interesting to experts on applying
of electron beam technologies for obtaining different types of nanostructured ceramics coatings.

**Keywords:** Physical Vapor Deposition, Reactive Evaporation, Nanostructurized Ceramics Coatings, Electron Beam Evaporation, Arc Discharge, High Voltage Glow Discharge, Residual Gas

**INTRODUCTION**

Today in the technologies of obtaining the nanostructured ceramic thin films is widely used the reactive deposition. The main physical processes of reactive films deposition are heating of evaporated metals by the powerful electron beam and propagation of metals vapor in the medium of active gases with lighting and maintaining the arc discharge. The distinguishing feature of such evaporation process is that in ionized plasma of arc discharge can take place such specific chemical reactions, which maintaining in the ordinary conditions, without gas ionization, is impossible. By applying such advanced electron beam technologies of films deposition can be obtained the perfect nanostructurized ceramic coatings with unique properties. Obtaining of such kind of chemical compound at the ordinary conditions, without additional ionization of metal vapors and technological gases in the arc discharge, is impossible. For example, obtaining of titanium nitride and titanium carbide films possible only at the conditions of reactive evaporation [1 – 4].

Since the traditional electron guns with heated cathodes can’t long time stably operated at the medium of active gases, often for providing the reactive deposition of ceramic thin films applied the high voltage glow discharge electron guns (HVGDEG). Generally, such guns are always characterized by the simplicity of construction, cheapness of the gun and evacuation equipment, as well as by the stable operation in the soft vacuum, range of 1 – 10 Pa, at the medium of active gases [5, 6]. In additional, effective control of beam power both by changing the operation pressure at the gun chamber [7] and by lighting the additional low-voltage discharge [8] is also possible.

Advanced technology of obtaining nanostructurized ceramic coatings with using HVGDEG can be successfully applied for many industrial applications in mechanical engineering, automotive industry, aerospace industry, as well as instrument-making and in electronics industry. The ceramic films, obtained with applying reactive deposition by using HVGDEG as the source of electron beam, are usually characterized by the high mechanical strength, heat resistance, as well as by the good dielectric properties. Therefore, using of such ceramics films and coatings is very perspective at cutting instruments, as thermoinsulation coating in engines, as well as thin dielectric films in the modern electronic devices, including microwave devices for communication systems [9–11].

Unfortunately, further development of advanced technology of reactive deposition of thin films with using HVGDEG is hampered by the lack of precision and adequate mathematical models of this process. The absence of such kind mathematical model does not allow to elaborate effectively the electron beam technological equipment for reactive deposition of nanostructurized thin films. Since the basic principles of simulation of high voltage glow discharge (HVGD) electron sources generally have been considered at the paper [6], for forming the complex model of process of reactive deposition of thin films considering the basic principles of simulation the arc discharge in the metal vapors over the crucible is necessary. Therefore, the aim of this chapter is considering these basic physical principles for such kind of model, obtaining necessary analytical relations and analyzing obtained simulation results.

**BASIC STATEMENT OF CONSIDERED PROBLEM**

The generalized scheme of electrodes system for lighting and maintaining arc discharge in the metal vapors, which included surface of evaporated metal, located in crucible, and additional ring-like electrode with the positive potential relatively to crucible, is presented in Fig. 1 [5]. Basic geometry parameters of depicted at Fig. 1 electrodes’ system, which are necessary to forming the mathematic model, are presented at Fig. 2. On Fig. 2 the radius of cross-section of the ring electrode is noted as \( r \), the inner radius of ring – as \( R \), the angle between axis \( z \) and the line, which connected the origin of coordinate system with the center point of ring – as \( \alpha \), and the distance between electrode and crucible as \( d \). Corresponding to Fig. 2, the solving of considered task was provided in cylindrical coordinates system, and the origin of coordinates system is coincided with the center of upper base of cylindrical surface of crucible.
The mathematic model for simulation of electrodes system, presented schematically at the Fig. 2, is generally based on the following presumptions.

1. Distribution of electric field between the ring-like electrode and the surface of evaporated metal is defined by analytical solving of Poisson Equation (PE) at the cylindrical coordinates.

2. The value of space charge, which stand at the right side of PE, is defined by the equation of continuity for electrons and ions current.

3. Energy of the metal’s atom at the upper base of crucible surface defined from the Boltzmann law as $kT_{ev}$.

4. The vapor under the crucible can be considered as saturated. Corresponding to this assumption defined the concentration metal’s atoms $n_{m}$.

5. Concentration of atoms of residual gas is defined by pressure in technological chamber $p_{g}$.

6. Level of ionization of metal vapor atoms is defined by the coefficient of ionization $\beta_{m}$.

7. Level of ionization of residual gas atoms is defined by the coefficient of ionization $\beta_{g}$.

8. Corresponding to the model of single charged ions, the number of ions and electrons at the volume of arc discharge is considered as equal.

FORMING AND SOLVING OF BASIC EQUATION SYSTEM

Equations for calculation the distribution of electric field

Generally, the equation for defining of electric field at the electrodes system, which structure and geometry parameters are presented at Fig. 2, is follow from general form of Poisson equation at cylindrical coordinates:

$$\frac{d^2U(r,z)}{dr^2} = \rho.$$ (1)

where $U(r,z)$ – the electric potential, $\rho$ – space charge density [12, 13].

Analytical solving of equation (1) for electrodes system with the plane and thin-ring electrodes, like plotted at Fig. 2, is written as follows [12, 13]:

$$U(\alpha) = \frac{\rho \cdot \sin(\alpha)}{2\varepsilon \varepsilon_0} = \frac{Q \cdot \sin(\alpha)}{4\pi \varepsilon_0 R}, \quad Q = 2\pi R, \rho.$$ (2)

where $\varepsilon$ – dielectric constant of residual gas, $\varepsilon_0 = 8.85418 \times 10^{-12} F/m$ – universal dielectric constant, $Q$ – space charge of elementary ring [12, 13]. Corresponded to equation (2), axial strength of electric is defined as [12, 13]:

$$E(\alpha) = \frac{Q \cdot \sin^2(\alpha) \cos(\alpha)}{4\pi \varepsilon_0 R^2}.$$ (3)
Defining the pressure of saturated vapor, concentration of metals’ atoms and necessary power of electron beam

The pressure of saturated vapor at the system of reactive deposition of ceremic films, which general constructive scheme is presented at Fig. 1 and the geometry parameters are defined at Fig. 2, in the general case can be defined with using Mendeleev – Clapeyron equation as follows [14 – 16]:

\[ p_s = \frac{\rho_s}{\mu_v} RT_{ev}, \]  
(4)

where \( p_s \) – pressure of saturated vapors, \( \rho_s \) – density of vapors, \( \mu_v \) – molecular mass of vapors, \( R = 8.314 \frac{J}{mol \cdot K} \) – universal gas constant, \( T_{ev} \) – temperature of metals evaporation. Since the density of vapors \( \rho_v \) is always connected with molecular mass and pressure by the following relation [14 – 16]:

\[ \rho_v = \frac{\mu_v N_A p_s}{p_a}, \]  
(5)

where \( N_A = 6.022 \cdot 10^{23} \frac{1}{mol} \) – Avogadro constant, \( p_a \) – atmospheric pressure, the equation (4), taking into account (5), can be rewritten as follows:

\[ p_s = \frac{p_a T_{ev} N_A R}{p_a}. \]  
(6)

After mathematical transformers of the relation (6), the direct analytical expression for pressure of saturated vapors is simply written as follows:

\[ p_s = \sqrt{p_a T_{ev} N_A R}. \]  
(7)

In another hand, saturated pressure is connected with concentration of metals atoms \( n_m \) by the following analytical relation [14 – 16]:

\[ p_s = n_m kT_{ev}, \]  
(8)

where \( k = 1.380649 \cdot 10^{-23} \frac{J}{K} \) – Boltzmann constant. Therefore, taking into account relations (7) and (8), the analytical relation for concentration of the atoms of metal vapors over the crucible is written as follows:

\[ n_m = \frac{p_s}{kT_{ev}} = \frac{\sqrt{p_a T_{ev} N_A R}}{kT_{ev}} = \frac{1}{k} \sqrt{p_a T_{ev} N_A R}. \]  
(9)

Concentration of metal’s atoms \( n_{ev} \) have been defined with using the equation (9), is used later by substitution to the differential equation (1) for defining the space charge in simulated electrodes system of reactive deposition of thin films. Dependence \( n_{ev}(T_{ev}) \), have been obtained with using relation (9) is presented at Fig. 3.

It is also important, that concentration of vapor atoms \( n_{ev} \) is strongly depended on power of electron beam and defined by simple relation [14 – 16]:

\[ n_{ev} = \frac{P_{ev}}{q_{ev} m_w} = \frac{P_{ev} N_A}{q_{ev} m_w}, \]  
(10)

where \( q_{ev} \) – specific heat of vaporization, \( m_w \) – mass of metal atoms.

Taking into account relation (9), the dependence \( P_{ev}(T_{ev}) \) for obtaining the regime of saturation vapors, can be written as follows:

\[ P_{ev}(T_{ev}) = \frac{q_{ev} \mu_v}{k} \sqrt{\frac{p_a R}{T_{ev} N_A}}. \]  
(11)

Dependence \( P_{ev}(T_{ev}) \), obtained with using relation (11) for titanium evaporation which have thermodynamic parameter \( q_{ev} = 4.226 \cdot 10^5 \frac{J}{mol} \) and \( \mu_v = 47.88 \cdot 10^{-3} \frac{kg}{mol} \) [17], is presented at Fig. 4.
Defining the space charge in the system of reactive deposition of thin films on the base of equation of continuously the current of arc discharge

Let’s rewriting the equation (1) for considered electrodes system, applied for reactive deposition of ceramic coating, with taking into account relation for concentration of metals’ atoms (9) and other relations of gas dynamic [14 – 16]. Firstly, the current density of arc discharge has to be estimated on the base of fundamental relations of gas-discharge theory, presented in monographs [18 – 20]. Generally, the current density of arc discharge can be estimated as follows [18 – 20]:

\[
j_d = e n_{im} \left( \sqrt{\frac{2 k T_i}{m_i}} + \sqrt{\frac{2 \varphi(r)}{m_i}} \right) + e n_g \left( \sqrt{\frac{2 k T_0}{m_g}} + \sqrt{\frac{2 \varphi(r)}{m_g}} \right),
\]  

(12)

where \( m_i \) – mass of metal ions, \( m_g \) – mass of gas ions, \( n_{im} \) – concentration of metal ions, \( n_{ig} \) – concentration of gas ions, \( T_0 \) – the temperature of environment, \( n_e \) – concentration of electrons, \( m_e \) – mass of electrons.

For model of single-charge ionization, taking into account the law of charges equilibrium, such relation is always holds [18 – 20]:

\[
n_e = n_{im} + n_{ig}.  
\]

(13)

Taking into account, that concentration of gas atoms is defined as [14 – 16]:

\[
n_g = \frac{p_g}{k T_0}  
\]

(14)

to relation (13), with taking into account (9), (14), can be rewritten as:

\[
n_e = \beta_m \frac{p_w N_A R}{T_{ev}} + \beta_e \frac{p_g}{k T_0} = \frac{1}{k} \left( \beta_m \sqrt{\frac{p_w N_A R}{T_{ev}}} + \frac{\beta_e p_g}{T_0} \right),
\]

(15)

where \( \beta_m \) – level of ionization of metal vapor atoms, \( \beta_g \) – level of ionization of gas atoms.

Therefore, equation (12), with taking into account (9, 14, 15), can be rewriting as follows:

\[
j_d = \frac{e}{k} \left( \beta_m \frac{p_w N_A R}{T_{ev}} + \frac{2 k T_i}{m_i} + \frac{2 \varphi(r)}{m_i} + \beta_e \frac{p_g}{T_0} \right) \sqrt{\frac{2 k T_0}{m_g}} + \left( \beta_g \sqrt{\frac{2 k T_i}{m_g}} + \frac{2 \varphi(r)}{m_g} \right).
\]

(16)

For simplifying obtained equation (16) and further solving the differential equation (1), the following coefficient have been introduced:

\[
C_1 = \frac{e \beta_m}{k} \frac{p_w N_A R}{T_{ev}}; \quad C_2 = \frac{e \beta_e}{k T_0}; \quad C_3 = C_1 + C_2.
\]

(17)

With using substitutions (17), equation (16) has been rewritten in the simplified form as follows:

\[
j_d = C_1 \left( \sqrt{\frac{2 k T_i}{m_i}} + \frac{2 \varphi(r)}{m_i} \right) + C_2 \left( \sqrt{\frac{2 k T_0}{m_g}} + \frac{2 \varphi(r)}{m_g} \right) + C_3 \frac{2 \varphi(r)}{m_e}.
\]

(18)

With known the spatial distribution of electrical potential \( \varphi(r) \), the current-voltage characteristic of considered arc discharge at the metal’s vapors can be obtained with using equation (18).

Now, for further analytical solving of differential equation (1), the components of obtained equation (18) for current density of metal ions \( j_{im} \), of residual gas ions’ \( j_g \), and of electrons’ \( j_e \), have to be considered separately. In such case, the following analytical relations should to be written:
It is also well-known fact, that the space charge, formed by the moving charged particles in the electrodes' system, is defined form the current density of this particles by the following relations [12, 13, 18 – 20]:

\[
\rho_s = \frac{j_s}{v_s} = \frac{2q_e \varphi(r)}{m_e},
\]

(20)

where \( s \) – sort of the particles, \( v_s \) – velocity of the particles, \( q_e \) – charge of the particles.

Therefore, for obtaining the values of space charges, formed by the metals' ions \( \rho_{\text{m}} \), by the residual gas ions \( \rho_{\text{g}} \) and by the electrons \( \rho_s \) corresponding to equations (19, 20), such analytical relations can be written:

\[
\rho_{\text{m}} = C_1 \left( \sqrt{\frac{kT_{ce}}{\varphi(r)}} + 1 \right),
\]

\[
\rho_{\text{g}} = C_1 \left( \sqrt{\frac{kT_{ce}}{\varphi(r)}} + 1 \right),
\]

(21)

Substituting coefficients \( C_p, C_j \) and \( C_j \) from relations (17) to relations (21) and making corresponded analytical transformers, the following relations to the values of space charges of ions and electrons can be written as follows:

\[
\rho_{\text{m}} = \beta_m \sqrt{p_a N_A R} + C_1;
\]

\[
\rho_{\text{g}} = \beta_g p_g \frac{e}{kT_0 \varphi(r)} + C_2;
\]

\[
\rho_s = C_3.
\]

(22)

Taking into account the law of equilibrium of charged particles, defined by the relations (13, 17), from obtained relations (22), after corresponded mathematical transformers, the formula for summarized space charge \( \rho_s \) can be wrote as follows:

\[
\rho_s(r) = \rho_{\text{m}} + \rho_{\text{g}} - \rho_s = \left( \sqrt{\frac{e}{k}} \left( \beta_m \sqrt{p_a N_A R} + \frac{\beta_g p_g}{\sqrt{T_0}} \right) \right).
\]

(22)

where

\[
K = \frac{e}{k}\left( \beta_m \sqrt{p_a N_A R} + \frac{\beta_g p_g}{\sqrt{T_0}} \right).
\]

(23)

For further theoretical analysis, it is very important, that the coefficients of the model \( C_p \), \( C_j \) and \( C_j \) have been cancelled out and, thus, vanished from the set of equations (22, 23). Consequently, the values of these intermediate model coefficients, in any way, don't affect to the potential distribution in the simulated electrode system. From a physical point of view, this fact can be explained by the compensation of the space charges of electrons and ions, which generally corresponds to the system of equations (17).

On the base of obtained relations (22, 23) the potential distribution in the considered electrodes system, which basic construction and main geometry parameters are presented at Fig. 1 and Fig. 2, can be defined by analytical solving the differential equation (1).

**Defining the distribution of electric field**

Taking into account the simple analytical relation (22), differential equation (1) can be written for the considered system of reactive deposition of thin films as follows:

\[
\frac{d^2 \varphi(r)}{dr^2} = K \frac{1}{\sqrt{\varphi(r)}},
\]

(24)

where the constant \( K \) defined by the obtained analytical relation (23).

The initials conditions for Cauchy problem [21] for formulated task, corresponded to Fig. 2, can be written as follows:

\[
\varphi(R_i) = U_i; \quad \frac{d \varphi(r)}{dr} \bigg|_{r=R_i} = 0.
\]

(25)

Corresponded explanation draft is given at Fig. 5.

![Fig. 5: Basic principle for defining the initial conditions for equation (24), corresponding to Fig. 2](image)

Therefore, solution of equation (24), with taking into account the initial conditions (25), is follows:

\[
\frac{d \varphi(r)}{dr} = K \int \frac{dr}{\sqrt{\varphi(r)}} + E_0 = 2K \sqrt{\varphi(r)} + E_0;
\]

\[
\varphi(r) = \int \left( 2K \sqrt{\varphi(r)} + E_0 \right) dr,
\]

(27)
where $E_0$ and $U_d$ – the constants, which are defined by initial conditions (25).

From the equations (25 – 27), such analytical relation to the value $E_0$ can be written:

$$E_0 = -2K \sqrt{U_d}.$$  \hspace{1cm} (28)

Taking into account obtained relation (28) and substitution the value of $E_0$ into equation (27), following analytical relation can be obtained:

$$\frac{4K}{3} \varphi^{1.5}(r) + E_0 r + U_0 = U_d. \hspace{1cm} (29)$$

From obtained equation (29) the direct analytical relation for defining unknown value $U_d$ from the basic discharge electric and geometry parameters is written as follows:

$$U_0 = U_d + 2K \left( R_c \sqrt{U_d} - \frac{2}{3} U_d^{1.5} \right) \hspace{1cm} (30)$$

Taking into account obtaining relations (28, 30) for the coefficients $E_0$ and $U_0$ finally analytical relation (27) for the potential distribution $\varphi(r)$ can be rewritten as follows:

$$\varphi(r) = \frac{4K}{3} \varphi^{1.5}(r) - 2K \sqrt{U_d} r + U_d + 2K \left( R_c \sqrt{U_d} - \frac{2}{3} U_d^{1.5} \right). \hspace{1cm} (31)$$

Let assume in the equation (31) the following substitute:

$$\varphi(r) = t^2. \hspace{1cm} (32)$$

After that substitution, formula (31) can be considered as standard cubic equation with the following coefficients:

$$at^3 + bt^2 + d = 0, a = \frac{4K}{3}; b = -1;$$

$$d = -2K \left( (R_c - r) \sqrt{U_d} - \frac{2}{3} U_d^{1.5} \right) \hspace{1cm} (33)$$

Analytical solving of the set of equations (33) with using well-known Cordano formulas [22] give such result:

$$q = \frac{27}{32K^3} - 0,25 \left( 2U_d^{1.5} - 3(R_c - r) \sqrt{U_d} \right);$$

$$D = \left( \frac{24K^3 \left( \frac{2}{3} U_d^{1.5} - (R_c - r) \sqrt{U_d} \right) + 27}{2} \right)^2 - \frac{1}{4096K^6};$$

$$u = \sqrt{\frac{q}{2} + \sqrt{D}}; v = \sqrt{\frac{q}{2} - \sqrt{D}}; y = u + v; t = \frac{1}{4K \sqrt{I}}.$$  \hspace{1cm} (34)

The simulation results for distribution of potential at the plane of ring electrode location, obtained with using relation (34), as well as results of calculation of current density distribution, will be present at the next part of this chapter.

**Simulation results and its analyze**

The results of simulation of distribution the electric field $U(r)$ at the plane of symmetry of ring electrode of simulated electrodes system, designed for reactive deposition of thin films at the soft vacuum, is presented at Fig. 6. These results have been obtained with using the set of equations (23, 34) for different values of arc discharge voltage $U_d$. These results was obtained for such discharge parameters:

1. The inner radius of ring electrode $R_c = 0.05$ m.
2. Power of electron beam $P_b = 20$ kW.
3. Pressure of residual gas at the technological chamber $p_g = 5$ Pa.
4. Level of ionization of metal vapor atoms $\beta_m = 0.8$.
5. Level of ionization of residual gas atoms $\beta_g = 0.75$.

![Graph showing distribution of electric field at the symmetry plane of ring electrode for the simulated system of reactive deposition of thin films with maintaining arc discharge.](image)

**Fig. 6:** Distribution of electric field at the symmetry plane of ring electrode for the simulated system of reactive deposition of thin films with maintaining arc discharge

Calculation was provided for evaporation of titanium at the medium of nitrogen, necessary empirical coefficients was taking from the references [17 – 21]. It is clear generally, that obtaining the precision values for level of gas and vapors...
ionization, $\beta_g$ and $\beta_m$, is also very sophisticated task. Usually level of ionization for different metals and gases is strongly depend on electron beam power, pressure in technological chamber and voltage of arc discharge maintaining [18 – 20]. In any case, obtained simulation results for distribution of electric field at the ring plane are looks plausible. Near the ring the potential of electric field is equal to value $U_d$ and the derivation of potential $\frac{dU(r)}{dr} \bigg|_{r=Rr} = 0$, that fully corresponded to boundary conditions (25). Neary the symmetry axis of simulated electrode system the electric potential is always few Volts smaller, and generally the reducing of potential can be explained by the sagging of electric field at the ring plane. This physical effect is also described mathematically with using equations (2, 3) [12, 13]. For free space, without charged particles, and for the values of voltage at the ring electrode $U_d$ range of 50 – 100 V, the sagging of electric field potential is generally grater, up to 15 – 20 V [12, 13]. Therefore, in the considered conditions of maintaining the arc discharge the smallest value of potential sagging near the axis of electrodes' system can be explained by the influence of space charge of the positive ions of metal vapors and residual gas [18 – 20].

As for space charge of electrons, its influence in the discharge systems is always smaller by the reason of small mass of electrons and its high velocity and negligible [18 – 20]. Also, it is important on the theoretical point of view, that derivation of electric potential on the axis of electrodes system is also

$$\frac{dU(r)}{dr} = 0,$$

This important condition is always satisfied for axial electrodes systems [12, 13].

With known dependence of potential distribution at the plane of symmetry of ring electrode the value of current density of arc discharge can be defined with using equation (18). Corresponded graphic dependences, obtained as results of using described simulation technique, are presented at Fig. 7, a, b. The values of the coefficients $C_1$ and $C_2$ for considered parameters of arc discharge, pointed out above, was

$$C_1 = 0.0111 \frac{C}{m^3}, C_2 = 0.0023 \frac{C}{m^3}.$$

Therefore, as it is clear form dependences (35, 36), the dependence of current density on electric potential $j_d(\phi(r))$ defined by the constantly grows mathematic function and namely by this reason the dependences $j_d(r)$, plotted at Fig. 7, and $U(r)$, plotted at Fig. 6, are generally similar.
As a result of provided researches, main analytical dependences for pervious estimation the electrical parameters of arc discharge at the technological systems of reactive deposition of thin films, have been obtained and described in this chapter of book. For real systems of films deposition, which are put into operation at the industrial equipment, more precision values of model coefficients $C_1, C_2, C_3, K, S_1, S_2$ can be defined empirically. By the such way experts can define the values of power of electron beam $P_o$, residual gas pressure $p_g$ and the voltage of arc discharge $U_o$, which are provided the suitable current density $j_d$ for maintaining the chemical reaction between the metal vapors and residual gases. As a result of applying the proposed model, the required quality of both standard and novel nanostructurized ceramic films, which can be obtained, is ensuring. For estimation the necessary power of electron beam, the relations (4 – 11), presented at the section 3.2, as well as graphic dependence, given at Fig. 4, can be used. For defining the necessary parameters of HVGDEG, which is generally considered as advanced tool for obtaining the nanostructurized thin ceramics’ films by using physical vapor deposition at the medium of active gases [1 – 4], can be used universal complex model, have been described early in the chapter of book [6]. And for defining the value of arc current density, necessary for maintaining the chemical reaction between the metal vapor and residual gas, the general recommendations, formulated in the papers [1 – 4], are usually used by the experts.

**CONCLUSION**

The mathematical model, based on the analytical relations, predetermined to pervious estimation the distribution of electrical potential and current density at the electron-beam technological systems for reactive chemical vapor deposition, is presented at this chapter of the book. At the second part of chapter the basic principles of formed model have been formulated and main electrical and geometry parameters of simulated arc-discharge electrodes’ system are defined.

At the section 3.1 formulated the basic set of equations for finding the electric potential distribution at the system with the positive potential on the thin ring electrode with inner radius $R_s$.

At the section 3.2 the main internal thermodynamic parameters of simulated evaporation system, such as pressure of saturated vapors $p_s$, temperature of metal evaporation $T_{ev}$ and the pressure of residual gas $p_g$ are considered and analytical relations for calculation the concentration of metals atoms $n_m$ and the necessary electron beam power $P_o$ have been proposed. Corresponded graphic dependences $n_m(T_{ev})$ and $P_o(T_{ev})$ are given at Fig. 3 and Fig. 4.

At the section 3.3 the basic mathematical relations for calculation the arc discharge current density and the space charge at the simulated electrodes’ system have been obtained, at this section of chapter the main analytical transformers, which have been used, are also shown.

At the section 3.4 the electrostatic problem for considered electrodes’ system has been solved by direct integration of Poisson equation at the cylindrical coordinates. The relations for space charge of moving particles, which are always necessary for solving the Poisson equation, have been obtained previously at the section 3.3. Direct integration of the Poisson equation led to obtaining the cubic equation, which have been solved analytically. At the section 3.4 also the main analytical transformers, which have been used, are shown.

At the section 4 the simulation results for the distribution of the electric potential and arc current density at the plain of ring electrode symmetry, obtained with using the presumptions and relations, given at the pervious parts of the chapter, have been presented and analyzed. Graphic dependences $U(r)$ and $j_d(r)$, for different voltage of arc discharge maintaining $U_o$ are given at Fig. 6 and Fig. 7.

The main advantage of proposed model of arc discharge for reactive deposition of nanostructurized ceramic films is its simplicity, because it based only on analytic relations and don’t include iterative numerical calculations. But in any case, it generally corresponded to the main physical principles of electric fields theory [12, 13], thermodynamic and gas dynamic theory [14, 15], as well as to the theory of gas discharges [18 – 20]. Therefore, generally obtained simulation results for distribution of electric field and arc discharge current density at the ring plane are looks plausible. For improving the accuracy of proposed model more precision values of the used semiempirical coefficients $C_1, C_2, C_3, K, S_1, S_2$ can be defined empirically for the specific technology and for metals and gases, applied for obtaining the high-quality nanostructure ceramic films [1 – 4]. Therefore, another advantage of the proposed model is its versatility and the possibility of its using for different parameters of the electrode system and for any metals and gases. Certainly, in such case defining of precision values of the semiempirical coefficients is the separate complex and sophisticated problem. For obtaining the approximative values of these coefficients the
handbooks on basic physics [17], on physics of gas discharges [18 – 20], as well as on gas dynamics [14 – 16] can be used.

It also should be pointed out, that for realizing the technology of electron-beam deposition of nanostructurized films the HVGDEG can be successfully applied, because the electron guns of such type are reliably and stably operated at the medium of active gases [5 – 8]. For defining the HVGDEG energetic parameters the universal complex model, proposed at the book [6], can be used. In the such situation the proposed early model of HVGDEG and the model of the arc discharge in the metal vapor, complexly described in this chapter, have to be used simultaneously. Therefore, at the practical point of view, these two models can be realized as one integrated software complex.

Basic analytical relations and simulation results, have been considered and analyzed in this chapter, are also very interesting at the practical points of view for the experts at the branch of designing the modern electron beam installations for deposition the nanostructurized ceramic coatings. Furthermore, the proposed model is simple and can be easy used by the experts. Choosing the parameters of technological process of nanostructured ceramic films deposition is the specific topic of further scientific and applied engineering researches. Therefore, applying of proposed model to the estimation of electric parameters of arc discharge in the metal vapors and active gases can led the experts to the new advanced ideas on creating the high-quality films form nanostructurized ceramics.

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