Hydrogen Behavior in an Ultrafine-Grained Electrodeposited Pure Iron

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The behavior of hydrogen in an ultrafine-grained electrodeposited pure iron with Lankford (r) value larger than 7.0 was studied by small angle neutron scattering (SANS) and thermal desorption spectroscopy (TDS). Nano-sized inhomogeneity consisting of hydrogen bubble exists in the deposited specimen. The bubble size increases a little by 673 K annealing and then all the bubbles disappear after 973 K annealing. With rolling at room temperature (RT), the bubble size and number density are found to decrease, which must be caused by the change in the status of hydrogen location during plastic deformation. Crystal rotation and grain coalescence are revealed to occur after rolling deformation from electron backscattered diffraction (EBSD) results.

KEY WORDS: electrodeposited pure iron; small angle neutron scattering; hydrogen; annealing; cold rolling.

1. Introduction

Hydrogen is easily introduced into metals and alloys during electrodeposition process and normally has negative influence on the quality of the deposits.1,2) It is not easy to experimentally detect hydrogen because of its low sensitivity to external stimulations, high mobility and small solubility. Therefore, the understanding of hydrogen behavior including diffusion and trapping has been limited due to the lack of suitable monitoring tools.3) Thermal desorption spectroscopy (TDS) is usually used to determine the hydrogen desorption characteristics, i.e., trapped sites and total amounts of hydrogen over the entire specimen.4–6) Information about the size, shape, number and distribution of hydrogen cluster or bubble are still uncertain. Small angle neutron scattering (SANS) is an appropriate technique for measuring microstructure in a scale from several nanometers to micrometers for many materials and has advantage of high sensitivity to light elements like hydrogen as neutron interacts with nuclei.7–9) In the present work, SANS and TDS were employed to investigate the hydrogen behavior in an electrodeposited pure iron with annealing or cold rolling. It has been reported that this electrodeposited pure iron sheets show the Lankford (r) value larger than 7.0.10–13) The reason of such unusual anisotropic plastic deformation is also discussed.

2. Experimental Procedures

Ultra-fine crystalline pure iron sheet used in this study was produced through industrial process by TohoZinc Co. Ltd., with purity of more than 99.995 wt.% As is observed in Fig. 1, microstructure of the as-deposited sheet consists of columnar grains with the cross section diameter of smaller than 1 μm (see (a) and (d)). The uni-directionally oriented columnar grains exhibited extremely sharp <111>//ND fiber texture ((Figs. 1(b), 1(c)). TEM microstructure reported in ref.12 ((Fig. 1(d)) suggests the existence of stress field. Some sheets with 1.5 mm in thickness were annealed at 673 K or 973 K for 3.6 ks and other 1.1 mm thick sheets were cold rolled at RT by 10%, 20% or 40% reduction in area. Scanning electron microscope (SEM) of Hitachi 4300, equipped with a TSL electron backscatter diffraction (EBSD) system was used for observing the top surface and cross section of the as-deposited specimens before and after deformation, and operated at an accelerating voltage of 15 kV and a tilt angle of 70° with step size of 0.1 μm.

TDS of hydrogen was conducted for the deposited and deformed specimens with a heating rate of 8 K/s in a temperature range from RT to 1 073 K in a vacuum chamber, where 10 mm×10 mm×1 mm specimens were prepared.5) Specimen size for SANS measurement was 20 mm×20 mm×2 mm. SANS measurements were carried out with a small-angle neutron scattering spectrometer SANS-J-II equipped at the cold neutron source of JRR-3, JAEA. A magnetic field of 1T was applied to separate nuclear scattering component from magnetic scattering component. And the scattered intensity was recorded by a two-dimensional detector which could be positioned at distances of 2.5 m and 10 m from the specimen. The incident wave length was set
to be 0.656 nm, which provided the scattering vector \( q = \frac{4\pi \sin \theta}{\lambda} \), \( 2\theta \) is the scattering angle, with neutron wave length of \( \lambda \) range of \( 5 \times 10^{-3} \text{nm}^{-1} \leq q \leq 1.99 \text{nm}^{-1} \), corresponding to real structure sizes of the order of \( 3 \sim 1000 \text{nm} \). The calibration to absolute intensity was performed by measurement of the irradiated aluminum standard specimen.

### 3. Results and Discussion

#### 3.1. Results of TDS

Figure 2 shows hydrogen desorption profiles obtained from the TDS for the as-deposited and the 40% rolled specimen. There are three main peaks (1), (2) and (3) on the as-deposited specimen spectra at about 473 K, 723 K and 953 K, respectively. All peaks are found to shift left and the intensity around peak (2) becomes stronger while that around peak (3) weaken after 40% cold rolling. In addition, the peak (2) and (3) seems to show split. It indicates that hydrogen distribution in the specimen was changed by cold rolling. It is reported that the activation energies for hydrogen evolution from different trapping sites in pure iron increases in the order of grain boundaries, dislocations and microvoid, and that hydrogen is claimed to exist in the microvoids as a molecule. After cold rolling, the lowest hydrogen release peak (1) shifted toward left, and intensity of peak (2) on TDS profile increased, which may correspond to hydrogen release from dislocation. On the other hand, the peak (3) related to hydrogen release from microvoid becomes much smaller. As the relative intensity of peak (2) becomes higher while that of peak (3) decreases in the 40% rolled specimen. It suggests that hydrogen trap site changed from microvoid to dislocation. The reasons for peak split at peak (2) and (3) must stem from some transition stages but are not clear at the moment.

#### 3.2. Effect of Annealing on SANS Profile

The SANS intensity in a magnetic field is composed of an isotropic nuclear contribution and an anisotropic magnetic contribution. The nuclear and magnetic components can be separated from each other using the following formula,

\[
I(q) = I_{\text{Nuclear}} + \sin^2 \alpha \cdot I_{\text{Magnetic}}(q) \quad ................. \quad (1)
\]

Where, \( \alpha \) is the angle between the scattering vector \( q \) and the magnetic field direction. The nuclear component of the scattering intensity can be determined at \( \alpha = 0^\circ \), while the sum of nuclear and magnetic intensity is found at \( \alpha = 90^\circ \).

**Figure 3** shows SANS profiles for the annealed specimens with a magnetic field of 1T. As seen, scattering intensity decreases with increasing of annealing temperature. So called “Guinier region” caused by the existence of nano-size inhomogeneity which causes SANS scattering contrast, for example, hydrogen bubble, is evidently observed in nuclear scattering profiles for the as-deposited and the 673 K annealed specimens, while not for the 973 K specimen. This indicates the disappearance of inhomogeneity by annealing at 973 K. It is interesting to compare this result with \( r \)-values reported in ref.12 (7.6 for the as-deposited, 5.2 for the 673 K annealed and 2.2 for the 973 K annealed specimen). The average particle radius determined from the fitting with
Guinier approximation\textsuperscript{20,21} were about 15 nm commonly for the as-deposited and 673 K annealed specimens. A remarkable change in SANS profile for the 973 K annealed specimen may be ascribed to a few reasons, i.e., defects like void, impurity elements like C and N, and hydrogen. By annealing at 973 K, the microstructure changes from fine columnar grains to coarse equiaxed grains as was reported in ref.\textsuperscript{12}. This change in grain size and shape is postulated to affect the SANS intensity at a lower $q$ region than that in the measured region. If carbides or nitrides precipitate by annealing, this would increase the SANS intensity. Hence, the drastic decrease in the SANS intensity observed at $q = 0.02–1$ nm\textsuperscript{–1} in Fig. 3 is believed to be caused mostly by the hydrogen desorption from comparison between Figs. 2 and 3.

The following equation can be used for profile fitting to get the particle size distribution\textsuperscript{19–21} assuming spherical inhomogeneity (independent spherical particles model).

\begin{equation}
I(q) = \Delta \rho^2 N_s \int_0^\infty V^2(R) F^2(q, R) N(R) dR \quad \ldots \ldots \ldots (2)
\end{equation}

In which, $\Delta \rho$ is the scattering length contrast between the particle and the matrix, $F(q, R)$ is particle shape factor, $N_s$ is number density of the particle, $F(R)$ is the volume of the particle. For spherical particle, the shape factor $F(q, R)$ can be written as,

\begin{equation}
F(q, R) = \left( \frac{3 \sin(qR) - qR \cos(qR)}{(qR)} \right) \quad \ldots \ldots \ldots (3)
\end{equation}

Eq. 4 depicts the log-normal normalised size distribution function,

\begin{equation}
N(R) = \frac{1}{R \sigma \sqrt{2\pi}} \exp \left( -\frac{\ln(R - \ln R_m)^2}{2\sigma^2} \right) \quad \ldots \ldots \ldots (4)
\end{equation}

Where $R$ is the particle radius, $R_m$ is the median radius and $\sigma$ is the variance of the distribution. After fitting, particle information including average size ($R_{ave}$), size distribution ($N(R)$), number density ($N_s$), the volume of the particle and volume fraction ($V_f$) can be calculated by the following equations.

\begin{equation}
R_{ave} = \frac{\int_0^\infty N(R)RdR}{\int_0^\infty N(R)dR} \quad \ldots \ldots \ldots (5)
\end{equation}

\begin{equation}
V_f = \frac{4}{3} \pi R^3 N(R)dR \quad \ldots \ldots \ldots (6)
\end{equation}

\begin{equation}
V_f = N_s \ast V_r \quad \ldots \ldots \ldots (7)
\end{equation}

The fitted profiles for the deposited and the 673 K annealed specimens were depicted in Fig. 4 showing good coincidence with the measurements, in which bubble quantity was computed based on the difference of scattering length density contrast ($\Delta \rho$) between hydrogen bubble assuming 1atm (0\textdegree C, 101.325 kPa) and the iron matrix. Scattering length density is defined as:

\begin{equation}
\rho = \sum_{i=1}^{n} \frac{b_{ci}}{V_m} \quad \ldots \ldots \ldots (8)
\end{equation}

\begin{equation}
\frac{1}{V_m} = N_A \ast \rho_0 \ast \frac{A}{A} \quad \ldots \ldots \ldots (9)
\end{equation}

Where $b_{ci}$ is the bound coherent scattering length of ith of n atoms in a molecular with molecular volume $V_m$, $N_A$ is Avogadro constant, $\rho_0$ is density of molecule and $A$ means relative atomic or molecule mass.

For hydrogen molecule, $\rho_0 = 0.08988$ g/L, $A = 2.016$ g/mol\textsuperscript{–1}, $b = (-3.739 \times 10^{-13} \text{ cm}) \times 2$, while for iron, $\rho_0 = 7.874 \text{ g/cm}^3$, $A = 55.845$ g/mol\textsuperscript{–1}, $b = 9.45 \times 10^{-13} \text{ cm}$, the calculated scattering length density is $\rho_{H_2} = -2.02 \times 10^7 \text{ cm}^{-2}$ and $\rho_{Fe} = 8.02 \times 10^{10} \text{ cm}^{-2}$, respectively.

\begin{equation}
\Delta \rho = \rho_{H_2} - \rho_{Fe} \quad \ldots \ldots \ldots (10)
\end{equation}

As is shown in Table 1, from SANS results, hydrogen bubble size became larger after 673 K annealing. However, bubble number density and volume fraction decreased because of hydrogen desorption. The hydrogen volume fraction in the as-deposited specimen estimated by SANS was in excellent agreement with that measured by TDS.

\subsection*{3.3. Effect of Cold Rolling Deformation on SANS Profile}

1 Tesla magnetic field was applied in the horizontal direction during SANS measurement on the cold rolled specimens. As seen in Fig. 5, the nuclear scattering intensity decreases obviously with increase of rolling reduction, and “Guinier region” is observed in all profiles. Fig. 5(a) shows the measured and fitted nuclear scattering profiles for the
cold rolled specimens with different rolling reduction. Similarly to the case of annealing described above, corresponding hydrogen bubble information was calculated and the results were presented in Figs. 5(b) and 5(c). It is found that average size, number density and volume fraction are decreased with increase of rolling reduction. Table 1 compared hydrogen content of the as-deposited and the 40% rolled specimens determined from SANS and TDS results. For the as-deposited specimen, the SANS result was consistent with that of TDS. However, the amounts of hydrogen in the 40% rolled specimen determined by SANS data was lower than that from TDS. This is because only hydrogen bubbles were considered in SANS, while all kinds of hydrogen forms including bubble, segregation and so on can be detected and transformed into hydrogen content by TDS. From the amount difference and TDS hydrogen desorption peaks changes observed in Fig. 2, some hydrogen bubbles are likely to transform into other existing forms by cold rolling.

3.4. Microstructure Change with Cold Rolling

Figure 6 compares the inverse pole figure map on the surface and corresponding ODF texture calculated from EBSD data. After 40% rolling, texture component in ND direction hardly changes, while that along the rolling direction the $<112>$ and $<110>$ orientations texture becomes

### Table 1. Influences of annealing or cold rolling on size, number density and volume fraction of inhomogeneity (hydrogen) measured by SANS and TDS.

| Specimen           | SANS by a sphere fitted method | TDS          |
|--------------------|--------------------------------|--------------|
|                    | Rave/nm | $Vr$/nm$^{-3}$ | Nd/cm$^{-3}$ | Vf/%  | Wt./ppm |
| Annealed Deposited | 3.36654 | 470.694 | 1.74906e+16 | 0.823 | 0.09  | 0.788  |
| 673 K annealed     | 4.49462 | 806.504 | 1.49534e+15 | 0.121 | –     | –      |
| Cold Rolled Deposited | 4.95829 | 1 024.59 | 1.77688e+16 | 1.82  | 0.09  | 0.788  |
| 10% rolled         | 3.53605 | 540.134 | 1.50534e+16 | 0.813 | –     | –      |
| 20% rolled         | 2.76265 | 382.475 | 9.23594e+15 | 0.353 | –     | –      |
| 40% rolled         | 2.67044 | 357.75  | 2.48938e+15 | 0.089 | 0.11  | 0.963  |
stronger. That means strong \{111\} <hkl> random texture evolved into \{111\} <112> and \{111\} <110> orientation dominant due to lattice rotation. All boundaries are defined as either low-angle grain boundaries with misorientation angle $\theta$ between 2° and 15° or high-angle grain boundaries above 15°, grain coalescence is verified from the EBSD grain boundary map on the surface by OIM software shown in Figs. 7(a) and 7(b).\textsuperscript{22–29} It may come from grain boundary motion by a kind of hydrogen related stress-driven mechanism during deformation.\textsuperscript{25–28} And the fraction of low angle boundaries increased from about 28% in the as-deposited specimen. Grain coalescence of the columnar as-deposited specimen. Grain coalescence of the columnar is recognized by comparing Figs. 1(a) and 7(c). As will be reported elsewhere, crystal rotation and grain coalescence are also observed in case of tensile deformation.\textsuperscript{29} Hence, grain coalescence related with hydrogen behavior during plastic deformation is speculated to realize unusual high r-value.

4. Conclusions

Effect of annealing and cold rolling on hydrogen behavior in an electrodeposited pure iron sheet was studied by means of TDS and SANS. The obtained results would be summarized as follows.

(1) Three hydrogen desorption peaks appear on the TDS profile during annealing, at about 473 K, 723 K and 973 K for the as-deposited specimen. All peaks shift left and two peaks at higher temperatures merge into unity after 40% cold rolling, which should be caused by the change in the status of hydrogen during the plastic deformation.

(2) From SANS results, nano-sized inhomogeneity is found to exist in the as-deposited and 673 K annealing specimen. Its size increases a little by 673 K annealing and then disappear after 973 K annealing. By using an independent spherical particles model for the SANS data obtained, the hydrogen volume fraction in the as-deposited specimen was estimated was in excellent agreement with that measured by TDS.

(3) With rolling at RT, the bubble size and volume fraction are found to decrease. Compared with TDS results, some hydrogen bubbles are likely to transform into other existing forms by cold rolling.

(4) Strong random texture \{111\} <hkl> grows into \{111\} <112> and \{111\} <110> orientation dominant after cold 40% rolling. Meanwhile, grain coalescence occurs during plastic deformation, which could be related to the extremely high r value.

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