JUPITER AS AN EXOPLANET: UV TO NIR TRANSMISSION SPECTRUM REVEALS HAZES, A Na LAYER, AND POSSIBLY STRATOSPHERIC H2O-ICE CLOUDS

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ABSTRACT

Currently, the analysis of transmission spectra is the most successful technique to probe the chemical composition of exoplanet atmospheres. However, the accuracy of these measurements is constrained by observational limitations and the diversity of possible atmospheric compositions. Here, we show the UV–VIS–IR transmission spectrum of Jupiter as if it were a transiting exoplanet, obtained by observing one of its satellites, Ganymede, while passing through Jupiter’s shadow, i.e., during a solar eclipse from Ganymede. The spectrum shows strong extinction due to the presence of clouds (aerosols) and haze in the atmosphere and strong absorption features from CH4. More interestingly, the comparison with radiative transfer models reveals a spectral signature, which we attribute here to a Jupiter stratospheric layer of crystalline H2O ice. The atomic transitions of Na are also present. These results are relevant for the modeling and interpretation of giant transiting exoplanets. They also open a new technique to explore the atmospheric composition of the upper layers of Jupiter’s atmosphere.

Key words: eclipses – planets and satellites: atmospheres – planets and satellites: gaseous planets – planets and satellites: individual (Jupiter, Ganymede)

1. INTRODUCTION

Over the past two decades, more than 1800 exoplanets have been discovered (Borucki et al. 2010), approximately 65% of which are transiting. For a small sample of these planets—those which orbit bright stars and have large planet-to-star area ratios—their atmospheres may be explored through transmission spectroscopy. During the transit of a planet in front of a star, the stellar flux is partially blocked, but a very small fraction of the stellar flux, 10⁻⁴ for a Jupiter-like planet and a Sun-like star system, passes through the thin planetary atmosphere (if an atmospheric thickness of 0.1 times the radius of Jupiter is considered). Transmission spectroscopy has allowed the detection of atomic Na, H, He, C, N, O, H2O, CH4, and CO2 (Bellucci et al. 2004; Ehrenreich et al. 2006; Borucki et al. 2010; Tinetti & Deroo 2010; Sing et al. 2011; Huitson & Sing 2013). These observations, however, push the detection capabilities of space- and ground-based observatories to the limit, and the results are often a source of discrepancies in the literature (Ehrenreich et al. 2006; Gibson et al. 2011; Sing et al. 2011), which need to be addressed. Thus, observing planetary transits in our own solar system can serve as an invaluable benchmark and can provide crucial information for future exoplanet characterizations.

Here, we report the transmission spectrum of Jupiter, with a high signal-to-noise ratio (S/N), as if it were a transiting planet. Our technique is to observe Ganymede, which is in synchronous rotation around Jupiter, when crossing Jupiter’s shadow. During the eclipse, the spectral features of the Jovian atmosphere are imprinted in the sunlight that, after passing through Jupiter’s planetary limb, is reflected from Ganymede toward the Earth (see Figure 1). The ratio spectrum of Ganymede before and during the eclipse removes the spectral features of the Sun, the local telluric atmosphere on top of the telescopes, and the spectral albedo of Ganymede. Ganymede and Europa are practically atmosphere-less bodies and do not introduce any significant variability in the spectra. Similar observations have previously been applied to retrieve the Earth’s transmission spectrum through lunar eclipse observations (Palle et al. 2009).

2. OBSERVATIONS

Initially, we observed an eclipse of Ganymede on 2012 October 6 using LIRIS (Manchado et al. 2004) at WHT in La Palma Observatory, Spain. This experiment was later repeated, and the results were confirmed by observing a second eclipse with XSHOOTER (Vernet et al. 2011) at the Very Large Telescope (VLT) in Paranal Observatory on 2012 November 18. Here, we focus the discussion on the VLT data due to their higher S/N, but a detailed analysis of the WHT observations leads us to virtually identical results.

The larger aperture of VLT allowed us to take rapid measurements—these eclipses are only observable from the ground typically a few of times per year and the suitable observing window lasts several minutes—with high spectral resolution. At the same time, we extended our measurements into the visible and the near-ultraviolet regions, covering from 300 to 2500 nm in a single exposure. This is possible through a dichroic splitting of the beam in three arms: UV, VIS, and near-IR. The 0′5, 0′4, and 0′4 slit widths were used for each range, providing averaged R = 9100, 17,400, and 11,300, respectively. The observing method consisted of taking uninterrupted spectra of Ganymede in stare mode during its translation through Jupiter’s shadow. The telescope active optics was used for guiding at Ganymede’s non-sidereal rate in order to keep
the target properly centered within the slit width during the dark phases of the event. The apparent magnitude of Ganymede is \( V = 4.6 \), but it decreases by about eight magnitudes during the umbra phase (Smith 1975; Tsumura et al. 2014).

Observations started at 3:30 UT, with a seeing below 1″ that remained stable during the night. The telescope flexure compensation procedure was carried out at 04:00 UT in order to keep the three slits staring at the satellite during the darker phases. The penumbral phase started at 04:19 UT (see Figure 1), in which the planet blocks part, but not all, the direct sunlight. During this phase, the amount of direct sunlight reaching the satellite decreases with the progress of the eclipse, whereas the light refracted by the planetary atmosphere reaching the satellite increases. The eclipse was total when the umbra phase started, and all direct sunlight was blocked, at 04:40 UT. During the umbra, only refracted light from the planetary atmosphere reaches the satellite. The final out-of-the-eclipse phase was not observed due to the occultation of the satellite by Jupiter’s disk, which happened at 06:00 UTC.

The data reduction of the XSHOOTER was performed using the ESO pipeline Reflex. Each individual exposure had enough S/N to be analyzed independently. Nevertheless, the three CCDs have different readout times, and we took the VIS exposures as references. We then averaged the several near-IR spectra taken during a VIS exposure and the closest UV spectra in time to construct a full spectrum from 300 to 2500 \( \mu \)m. This way, 41 individual spectra were determined corresponding to sunlit Ganymede before the eclipse (B1), 8 during the penumbra (P1), and 3 during the umbra (U), at all wavelengths simultaneously.

The penumbra transmission spectra were determined by calculating the ratio between the spectra taken during P1 and the spectra taken just before the eclipse B1. The umbra transmission spectra were determined by calculating the ratio between the spectra taken during U and the spectra taken just before eclipse B1. Again, these ratios will cancel the telluric contribution of the local atmosphere, the solar spectrum intrinsic in the observations, and the spectral signatures of the satellite (Montañés-Rodríguez et al. 2006; Palle et al. 2009; Yan et al. 2014), leaving only the contribution from Jupiter’s atmosphere.

During these observations, the airmass decreases as the eclipse progresses (from 1.565 in B1 to 1.510 in the deepest U), making it easier to identify residual telluric features because they would be seen in emission in the (pen)umbra/bright ratios (see Figure 2).

### 3. MODEL SIMULATIONS

Transmission spectra of the occultation of the Sun through Jupiter’s atmosphere as observed from Ganymede have been computed, simulating the observations of the eclipse in the penumbra and within the first stages of the umbra. The extension (horizontal) of the Sun’s disk as it is setting on Jupiter’s horizon has been considered. Because of the strong refraction of Jupiter’s atmosphere, this makes it possible to sound Jupiter’s limb with moderate vertical resolution (a few tens of kilometers) at the lowest tangent heights.

The transmission spectra have been calculated by using the Karslruhe Optimized and Precise Radiative Transfer Algorithm (KOPRA; Stiller 2000). KOPRA is a well-tested line-by-line radiative transfer model that offers all the necessary physics for studying this problem. This code was originally developed for
Distribution of haze, CH$_4$, and H$_2$O ice in our model, together with geometrical penumbral spectra from Figure 2.

The extinction of the aerosol particles moving down to lower tangent heights.

Note how the sounded region shrinks as the eclipse progresses while of the water-ice particle cloud used in the computation of the transmission spectra. Altitude is taken as zero at the 1 bar pressure level. In previous spectra also in full sunlight illumination Ganymede spectra immediately before the penumbra phase, divided by a reference spectrum, and they start to show signs of extra absorption at 0.5889 and 0.5896 μm. In green are the two umbra spectra for which the S/N is good in this wavelength region.

**Figure 3.** (Top) Transmission spectra of Jupiter calculated for the early phase of the penumbra and at several stages during the umbra over the 0.5–2.5 μm spectral region. The spectra show the most prominent CH$_4$ bands, the extinction of the aerosol particles (haze), and the two distinct absorption peaks of water ice at 1.5 and 2.0 μm. Plotted with a gray line is one of our observed penumbral spectra from Figure 2(b). (Bottom) An illustration of the distribution of haze, CH$_4$, and H$_2$O ice in our model, together with geometrical information on the range of tangent altitudes covered in each of our simulations in the top panel, corresponding to the mean times when the (umbra)-penumbral spectra were taken. Altitude is taken as zero at the 1 bar pressure level. In addition, we have overplotted the profiles of the concentrations of aerosols and of the water-ice particle cloud used in the computation of the transmission spectra. Note how the sounded region shrinks as the eclipse progresses while moving down to lower tangent heights.

**Figure 4.** The evolution of Ganymede spectra during the observations centered on the Na I doublet wavelength region. The gray spectra are fully illuminated Ganymede spectra immediately before the penumbra phase, divided by a previous spectra also in full sunlight illumination (arbitrary reference spectrum). These spectra are basically featureless. In blue are the penumbra spectra divided by the same reference spectrum, and they start to show signs of extra absorption at 0.5889 and 0.5896 μm. In green are the two umbra spectra for which the S/N is good in this wavelength region.

**4. RESULTS**

In Figure 2(a), the penumbra spectrum of Jupiter is shown, representing a direct transmission component through the upper layers of the atmosphere of the planet. This is directly comparable to the information retrieved through a planetary transit. As expected, the spectra show the major absorption features of the most abundant component, CH$_4$. However, most interesting is the change in the spectral continuum: at shorter wavelengths (UV–VIS), transmission is lower due to the extinction caused by aerosols. Extinction by clouds and aerosol particles have a smooth wavelength-dependent extinction effect on the stellar flux that does not produce sharp features. Clouds and hazes have been tentatively detected in Hot Jupiter atmospheres using transmission spectroscopy from the Hubble Space Telescope (Ehrenreich et al. 2006; Pont et al. 2012). Our results give additional confidence to these findings.

In Figure 2(b), several umbra spectra of Jupiter are shown. As the eclipse progresses, the geometry allows us to sample lower tangent heights, and thus probe lower into Jupiter’s atmosphere. The planetary atmosphere also has the effect of refracting the stellar light (García-Muñoz et al. 2012). This refracted component becomes more important as the eclipse progresses deeper into the umbra. Thus, these spectra present a much more prominent absorption not only by the aerosol particles (Mie scattering) mainly in the near-IR but also by the use in Earth’s atmosphere and has been recently adapted to the atmospheres of Titan and Mars (García-Comas et al. 2011). Here, the reference Jupiter’s atmosphere, including pressure, temperature, and the species abundances, were taken from González et al. (2011). We include the major gaseous species, CH$_4$ and H$_2$, and other minor as C$_2$H$_2$, CO, H$_2$O, and NH$_3$. The concentration of NH$_3$ (only relevant in the lowest troposphere) was taken from Griffith et al. (1991). The only species with significant concentration that is not included is ethane (C$_2$H$_6$). The absorption of its strongest band in the region of 2.2–2.5 μm would be, however, masked by the stronger CH$_4$ bands in that region. The molecular spectroscopic data for all species have been taken from the HITRAN compilation, 2012 edition (Rothman et al. 2013).

Rayleigh scattering by molecular hydrogen and helium have also been taken into account by including the Rayleigh optical cross sections provided by Ford & Browne (1973) for H$_2$ and by Chan & Dalgarno (1965) for helium. Ro-vibrational absorption bands resulting from collisions between pairs of H$_2$–H$_2$ and H$_2$–He, the so-called collision-induced absorption or CIA, are also significant in the lower Jupiter atmosphere, where they form a smooth feature mainly in the 2.0–2.5 μm region. Our simulations include this absorption with absorption coefficients at low temperatures, derived by Borysow (2002) for H$_2$–H$_2$ pairs and by Borysow et al. (1989) and Borysow & Frommhold (1989) for H$_2$–He.

In addition, we also included in the simulations Mie scattering by water ice and aerosols. The optical properties of the crystalline water ice (real and imaginary part of the refractive index) were taken from Mastrapa et al. (2008) for a temperature of 150 K. The aerosol particles were assumed to have a mean radius of 0.25 μm in the range of 0.2–0.5 μm derived by Zhang et al. (2013) for the equatorial particles. The real part of the refractive index of these aerosol particles was taken from Khare et al. (2013) and the imaginary part from Zhang et al. (2013).
CH₄ absorbing bands. However, the most interesting feature is the possible detection of stratospheric H₂O ice cloud features at 1.5 and 2.0 μm.

In Jupiter, the visible cloud deck is composed of NH₃ ice condensation and spans from 700 up to 200 mbar at the equatorial region, near the tropopause, which is especially conspicuous in planetary images taken in the 890 nm strong methane band (West et al. 2004). Below the visible cloud layer, clouds formed by NH₃ ice particles, NH₄–SH solid particles, H₂O ice, and H₂ONH₃ liquid solution are also present, but they are hard to detect (Taylor et al. 2004). Above these levels, the major atmospheric constituent is CH₄, but recent analysis of observations made with the Cassini Imaging Science Subsystem indicates that the planet is wrapped up by haze at the lower stratosphere at the equatorial region and midlatitudes at ≈50 mbar, which rises at higher levels (p < 20 mbar) over the poles (Zhang et al. 2013). But the composition of this layer remains undetermined.

It is from these upper levels that our spectroscopic signals are coming from. Over the umbra evolution, H₂O-ice absorption features at 1.5 and 2.0 μm appear and later diminish (see Figure 2(b)) as the sounded tangent height on the planetary limb crosses the altitude (or pressure levels) where these particles are located. Their detection in our transmission spectra allows us to determine that Jupiter’s upper atmospheric hazes near the equator are at least partially composed of a cloud deck of H₂O-ice crystals. While these H₂O-ice absorption bands are well known, the spectral measurements of VIMS on board the Cassini spacecraft taken in 2000 did not detect them (Formisano et al. 2003).

In Figure 3(a), our model simulations of Jupiter’s umbra and penumbra transmission spectra are plotted. All major features of the measured spectra are simulated. The absorption spectral features at 1.5 and 2.0 μm can be very well reproduced in our model with crystalline water ice, needing a total column of about 10¹³ particles cm⁻² with a size of 0.01 μm located near the 0.5 mbar level, a water amount 420 times larger than that measured by Herschel (Cavalie et al. 2013), which found a water mass load of 1.5 × 10⁻⁷ g cm⁻² in the gas phase. While the good spectral match supports the identification of H₂O ice, the large amounts and the location at low pressure implied by the data are puzzling. Moreover, the presence of water ice at 0.5 mbar does not have an obvious explanation.

In terms of aerosol mass loading, Zhang et al. (2013) derived 1 × 10⁻⁶ g cm⁻² at low latitudes and 1 × 10⁻⁴ g cm⁻² at high latitudes. The H₂O-ice particle distribution we derive (log normal with sigma = 0.3 and r = 0.01 μm) gives an H₂O-ice mass load of 6.3 × 10⁻⁷ g cm⁻². Hence, although it is larger that their aerosol mass load for low latitudes, it is comparable to their mass load for high latitudes, using observations at 0.25 and 0.9 μm. Hence, there is no conflict between the derived H₂O-ice load and the results from previous literature. The haze parameters used in our model are in agreement with those recently reported for Jupiter’s equator by Zhang et al. (2013).

The solar atomic lines should disappear in the ratio between the (pen)umbra and bright spectra of Ganymede. However, some residual features remain, as several effects will alter the shape of the spectral lines, including (1) small Doppler shifts due to the Sun–Ganymede–Earth relative speeds, (2) changes in the solar region contributing to the spectra, (3) instrumental flexures, and (4) Raman scattering producing a center-to-limb brightening (Yan et al. 2014). This latter effect seems to dominate the spectral shape of the ratio spectra near atomic solar lines, where the line residuals have an inverted W-shape pattern.

In Figure 4, we plot several penumbra and umbra spectra (both divided by the direct sunlit Ganymedes spectrum) focusing on the region near the Na I doublet. Most of the lines are strongly deformed in a classical W pattern, as is the case of the Fe lines in between the Na I doublet and other solar lines at 586.2 and 591.4 μm. However, in the case of the Na I, there is a net absorption indicating the presence of a Na layer in Jupiter’s upper atmosphere. This Na I absorption starts to be clearly detectable in the last penumbra spectra, at high S/N, and in all umbra spectra (with progressively decreasing S/N). We plan to carry out a detailed study of these features in a future paper.

The presence of Na in Jupiter’s upper atmosphere can be explained by the deposition of either cometary impacts (Noll et al. 1995) or the continuous outward flux of Na from Io (Mendillo et al. 2004). The fact that our temporal series of sunlit Ganymede spectra do not show any signs of Na absorption, and neither do the first penumbra spectra, indicates that the origin of this absorption is located within Jupiter’s atmosphere, and not in the torus of Na trailing along Io’s orbit.

The WHT observations have not been discussed at length in this Letter; however, the results of that earlier campaign are similar to those obtained with VLT. In Figure 5, the transmission spectrum of Jupiter in one of the umbra of the WHT data is compared to one of the umbra obtained with the VLT. While there are differences in the spectra due to the fact that the eclipse geometry is not exactly the same and that WHT needs two different exposures (different times, different umbra depths) to cover the 0.9–2.5 μm range, they show essentially the same spectral signatures.

5. CONCLUSIONS

In summary, here we have determined the UV–VIS–IR transmission spectrum of Jupiter as if it were a transiting exoplanet. This transmission spectrum reveals the imprints of strong extinction due to the presence of clouds (aerosols) and hazes in Jupiter’s atmosphere and strong absorption features from CH₄. More interestingly, the comparison with radiative transfer models reveals a spectral signature, which is attributed here to a stratospheric layer of crystalline H₂O ice. The atomic transitions of Na are also present. These results are relevant for the modeling and interpretation of giant transiting exoplanets, but they also open a new technique to characterize the upper
layers of Jupiter’s atmosphere. Taking advantage of the scanning of faint features that this technique provides, observations of other satellite eclipses from space could help to set limits to the stratospheric water abundance in the upper layers of Jupiter’s atmosphere and provide a way to monitor the rate of cometary impacts on Jupiter (Smith et al. 1999), which, in turn, has consequences for the formation history of the solar system.

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