A Periodic Transmission Line Model for Body Channel Communication

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ABSTRACT Body channel communication (BCC) is a technique for data transmission exploiting the human body as communication channel. Even though it was pioneered about 25 years ago, the identification of a good electrical model behind its functioning is still an open research question. The proposed distributed model can then serve as a supporting tool for the design, allowing to enhance the performances of any BCC system. A novel finite periodic transmission line model was developed to describe the human body as transmission medium. According to this model, for the first time, the parasitic capacitance between the transmitter and the receiver is assumed to depend on their distance. The parameters related to the body and electrodes are acquired experimentally by fitting the bio-impedentiometric measurements, in the range of frequencies from 1 kHz to 1 MHz, obtaining a mean absolute error lower than 4° and 30° for the phase angle and impedance modulus, respectively. The proposed mathematical framework has been successfully validated by describing a ground-referred and low-complexity system called Live Wire, suitable as supporting tool for visually impaired people, and finding good agreement between the measured and the calculated data, marking a ±3% error for communication distances ranging from 20 to 150 cm. In this work we introduced a new circuital approach, for capacitive-coupling systems, based on finite periodic transmission line, capable to describe and model BCC systems allowing to optimize the performances of similar systems.

INDEX TERMS Body area network, body channel communication, intra-body communication, propagation model, periodic transmission line.

I. INTRODUCTION

The expressions intra-body communication, human body communication and Body channel communication (BCC) all refer to a short-range communication method that exploits the human body as a transmission medium for electrical signals [1], [2]. The pioneering work of Zimmerman on BCC was motivated by the dream and the imagination of a digital interconnected world where the center of communication is the human being [1], [2]. Driven by this idea, BCC can be established relying on two main coupling mechanisms [3]. The signal can be coupled in the body by direct current injection (galvanic coupling - GC) from a transmitter (TX) electrode to a receiver (RX). Alternatively, the near-field coupling of a quasi-electrostatic field between TX and RX electrodes can be used for (capacitive coupling - CC) [3]–[9]. However, the possibility of conveying information in the sub-cutaneous fat at frequencies around 2.45 GHz is currently
This human-centered interactivity is interesting because the communication is triggered by the person and its willingness to communicate with the external world [2], [3]. The nature of BCC is intrinsically safe, secure and resilient since the data remain confined within the biological medium, making this technology of great interest for applications such as e-payment, work-management, intelligent transportation system, security and entertainment. [23]–[26].

In order to refine and improve the performance of the BCC systems, it is fundamental to have a clear understanding of the fundamentals on-body CC communication. Indeed, despite the knowledge of the features and performances of the available devices, the understanding of the physics and modeling of the signal transmission and propagation in the biologic medium need to be improved. Therefore, this work aim to bridge the gap between the hardware design and the on-body propagation, by developing a model that can deal with the description of the complex mechanism of capacitive coupling, while including the hardware perspective, i.e., the TX and RX components.

In the literature it is possible to find several theoretical models, combined with circuital and/or full-wave (FW) approaches to the human channel modeling [3]–[6], [27]. Some models are capable to seize the essential features of the body channel propagation [28]. An example is the dipole model introduced by Bae et al. [8], in which the signal propagation in the CC depends on three distinct mechanism, namely surface wave, reactive radiation, and near-field quasi-static coupling [8], [9]. Regarding the BCC electric circuit model, both lumped and distributed parameter circuits have been investigated [25], [29]–[32]. The lumped parameter model can easily account for the essential electrical properties of different tissues, thus allowing to derive some closed-form expressions for the signal attenuation. However, the flaw of these models is their limited capability to represent realistic and complex cases [3]. Hence distributed models were also introduced. In particular, circuits diagrams were derived by cascading a basic RC cell, in order to analyze the influence of TX-RX distance on signal propagation [25], [33]. Although these simple models allow
deriving simple approximate expressions for the path loss, they are not able to describe the peculiar propagation mechanism of BCC [3]. Therefore, a transmission line model (TLM) using infinite structures was proposed by Callejón et al. [32] to study the dispersion and attenuation of the body channel considering the electrophysiological properties of the skin as a transverse admittance in the lumped model. However, in this model, the influence of the finite length of the channel on the communication performances is not analyzed in depth. Furthermore, the role and effects of the parasitic capacitance is marginal in the mathematical frameworks, despite its physical relevance [3]. As a result, the state of the art in the field lacks of mathematical models adequate to describe the physics of the on-body propagation using the TLM formalism, while accounting for the influence of the extrinsic parasitic return path on the communication.

Our work aims to bridge the gap for the case of CC mechanism. Indeed, the strategy of developing distributed circuitual models is preferable to the extensive, time-consuming and computational-demanding FW approach since the numerical simulations demonstrated that the electric field is almost entirely confined within the skin layer [8], [9], [34], and that the extrapolation of the parasitic capacitance is a non-trivial task [31].

In particular, in this work we developed a mathematical model capable of describing the body-channel propagation using CC devices. The proposed mathematical framework takes into account the strong influence of the environment, typical of capacitive devices, by taking into consideration the relevance of parasitic capacitances in the communication. The proposed distributed model can then serve as a supporting tool for the design, allowing to enhance the performances of any BCC system, especially in terms of working frequency choice, electrode sizing and signal strength.

In this work, the novel TLM model was experimentally validated on the signal propagation with Live Wire, a prototype of system which is currently used as indoor navigation aid for blind subjects [35].

The paper is organized as follows. The proposed body-propagation model based on finite periodic transmission line is explained in detail in Sect. II. The circuitual representation of CC system is presented in Sect. III. The use-case for the assessment of the model is introduced in Sect. III-A. The model parameters were retrieved from bio-impedentiometric measurements on a set of subjects. The experimental setup is presented in Sect. IV-A, along with the methodological information regarding the numerical fitting performed to derive the model parameters. The explanation of the transmission experiments and the channel characterization tests are reported in Sect. IV-B. The results of the impedance measurements, the fitting and the model validation are given in Sect. V-A. The findings related to the signal transmission using Live Wire in several configurations are provided in Sect. V-B. Finally, conclusions are drawn and presented in Sect. VI.

II. PROPAGATION MODEL

In order to describe, and analyze the propagation of the CC signal, it is possible to connect the physical properties of the on-body propagation channel to its electrical response. Assuming that the skin structure does not vary significantly along a body segment or region, the in- and on-body propagation can be described considering the propagation in a homogeneous structure along the direction of propagation [36], [37]. Since it is possible to derive lumped models for small skin parts, a TLM propagation model can be developed by considering the repetition of a sequence of identical cells. Therefore, the transmission line, having characteristic impedance $Z_0$, is assumed to be a periodic repetition of a lossy and dispersive RLC cell [37]. Therefore, the connection between the BCC physical mechanism and its electrical response can be done by modeling the channel as a finite periodic sequence of equal two-port networks, each one describing a sufficiently short section of the channel, as in Fig. 2. The complete circuit can be described using voltages ($V$) and currents ($I$) at the interfaces. Setting

$$W(q\Delta z) = \begin{bmatrix} V(q\Delta z) \\ I(q\Delta z) \end{bmatrix}$$

\hspace{1cm} (1)

where $\Delta z$ is the length of the section and $q = 1, 2, \ldots p$, being p the total number of repeating units. By using the transmission matrix of the unit cell $T$

$$T = \begin{bmatrix} A & B \\ C & D \end{bmatrix}$$

\hspace{1cm} (2)

we can write

$$W(p\Delta z) = T^p \cdot W(0)$$

\hspace{1cm} (3)

Eq. (3) can be given a closed-form expression using the eigenvalue decomposition of $T$, i.e., the eigenvalues $\lambda$ and the eigenvectors $h$, solutions of

$$T \cdot h = \lambda h$$

\hspace{1cm} (4)

The eigenvalues equation is

$$\lambda^2 - (A + D)\lambda + 1 = 0$$

\hspace{1cm} (5)

and therefore the product of the eigenvalues ($\lambda$) is equal to two. We can therefore let $\lambda_1 = e^{-jk\Delta z}$, $\lambda_2 = e^{jk\Delta z}$, where $k$ is a complex number, and obtain

$$\lambda_1 + \lambda_2 = e^{-jk\Delta z} + e^{jk\Delta z} = 2 \cos k \Delta z = A + D$$

\hspace{1cm} (6)

\footnote{For $q < 0$, we let, in Eq. (3), $T^q = T^{-|q|} = (T^{-1})^{|q|}$. Note that $T$ is always nonsingular since $det(T) = 1$.}
The eigenvectors can be normalized and therefore expressed as 
\[ h = \begin{bmatrix} 1 \\ y_1 \end{bmatrix}, \]
where
\[ y_1 = \frac{e^{-jk\Delta z} - A}{B} \]
\[ y_2 = \frac{e^{jk\Delta z} - A}{B} \]  

(7)

Using the eigenvectors we can represent \( T \) as
\[
T = H \cdot \begin{bmatrix} e^{-jk\Delta z} & 0 \\ 0 & e^{jk\Delta z} \end{bmatrix} H^{-1}
\]

(8)

where
\[
H = \begin{bmatrix} 1 & 1 \\ y_1 & y_2 \end{bmatrix}
\]

and
\[
H^{-1} = \frac{1}{y_2 - y_1} \begin{bmatrix} y_2 & -1 \\ -y_1 & 1 \end{bmatrix}
\]

obtaining
\[
T^q = H \cdot \begin{bmatrix} e^{-jkq\Delta z} & 0 \\ 0 & e^{jkq\Delta z} \end{bmatrix} H^{-1}
\]

(9)

From (9) it is possible to derive the voltages and the currents on the channel.

Since the unit cell is specular, i.e., symmetrical [38], \( A = D = \cos (k\Delta z) \) [see (6)], so that
\[ y_1 = -\frac{j \sin (k\Delta z)}{B} \]
\[ y_2 = \frac{j \sin (k\Delta z)}{B} = -y_1 \]  

(10)

and obtaining
\[
H^{-1} W(0) = -\frac{1}{2y_1} \begin{bmatrix} -y_1 & -1 \\ -y_1 & 1 \end{bmatrix} \begin{bmatrix} V(0) \\ I(0) \end{bmatrix}
\]
\[ = \frac{i}{2} \begin{bmatrix} 1 & 1/y_1 \\ 1 & -1/y_1 \end{bmatrix} \begin{bmatrix} V(0) + \frac{y_1}{y_1} I(0) \\ V(0) - \frac{1}{y_1} I(0) \end{bmatrix} = \begin{bmatrix} V^+ \\ V^- \end{bmatrix}
\]

(11)

where \( V^+ \) and \( V^- \) are the propagating and reflected voltage waves, respectively. It follows that
\[
W(z) = H \cdot \begin{bmatrix} V^+ e^{-j\kappa \Delta z} \\ V^- e^{j\kappa \Delta z} \end{bmatrix}
\]

(12)

i.e.,
\[
V(q\Delta z) = V^+ e^{-j\kappa q\Delta z} + V^- e^{j\kappa q\Delta z}
\]

\[ I(q\Delta z) = y_1 \left[ V^+ e^{-j\kappa q\Delta z} - V^- e^{j\kappa q\Delta z} \right]
\]

(13)

which is clearly analogous to the propagation along a transmission line.

**A. CHANNEL ANALYSIS**

The electrical response of the on-body channel depends on the parameters \( k \) and \( y_1 \) [see Eq. (11)], based on the physical behavior of the unit cell. The equivalent circuit representing the body channel is shown in Fig. 3.

Given the capacitive coupling mechanism and the low working frequency of BCC devices, the body propagation is assumed to be confined in the superficial skin layers (i.e., epidermis, dermis and subcutaneous layers) [3], [34]. On the skin, longitudinal and transverse paths can be recognized as relevant to the signal propagation, as underlined in [25], [32]. The propagation of the CC signal in the body channel is described by assuming the skin to be homogeneous along the propagation direction and, hence, the periodic transmission line formalism can be applied [36], [38]. The single cell considered in this study is shown in Fig. 3. The inductance \( L_b \) accounts for the radio-frequency currents induced in the body along the direction of propagation, whilst \( C_b \) accounts for the inter-electrode coupling and the parasitic effects of the environment, and \( R_b \) is the resistance of the body segment. The signal in the transverse section is supposed to flow as an electric current through the subject skin, via the skin capacitance \( C_s \), and then to the ground by a parasitic capacitance \( C_l \) between the skin and the external ground. The skin layer is modeled as a lossy capacitor with capacitance \( C_s \), in \( \text{Fm}^{-1} \) [25], [32], [39]–[41].

\[
C_s(\omega) = \varepsilon_\infty + \sum_{i=1}^{4} \frac{\Delta \varepsilon_i}{1 + \left( j\omega \tau_i \right)^{1-\alpha_i}} + \frac{\sigma}{j\omega \varepsilon_0}
\]

= \varepsilon_0 (\varepsilon'_s - j\varepsilon''_s)

(14)

where \( \varepsilon_\infty \) is the dielectric permittivity at optical frequencies, \( \Delta \varepsilon \) is defined as the difference between the static and the optical permittivities (i.e., \( \Delta \varepsilon = \varepsilon_\infty - \varepsilon_s \)), \( \tau \) is the relaxation time in \( s \), as found in [39]–[41] for the wet skin. The terms \( \varepsilon_0 \) is the vacuum dielectric permittivity in \( \text{Fm}^{-1} \) and \( \omega = 2\pi f \) is the angular frequency.
The transmission matrix of the circuit of Fig. 3 is [42]:

\[
T = \begin{bmatrix}
1 + (R + jX)(G + jB)\Delta z^2 & 2(R + jX)\Delta z + (R + jX)^2(G + jB)\Delta z^3 \\
(G + jB)\Delta z & 1 + (R + jX)(G + jB)\Delta z^2
\end{bmatrix}
\]  \hspace{1cm} (15)

Therefore, (6) becomes

\[
\cos (k\Delta z) = A = 1 + (R + jX)(G + jB)\Delta z^2  \hspace{1cm} (16)
\]

and, for negligible \( \Delta z \), a Taylor expansion of the right hand side gives the equation

\[
-\frac{1}{2} k^2 = (RG - XB) + j(RB + GX)  \hspace{1cm} (17)
\]

If the network was lossless, the equation for \( k \) would be \( k^2 = 2XB \), so that we have actual propagation for \( XB > 0 \). Since \( B > 0 \), we have propagation as long as the series LC in Fig. 3 has an inductive nature, i.e., \( X > 0 \). This is true in the frequency range of interest. However, as the frequency increases, the propagation would be cut-off and the signal strength would be attenuated when the capacitive susceptance becomes larger than the inductive one (see Fig. S2 of the Supp. Material). In the real lossy case, we have small losses \( (R \ll |X|, G \ll |B|) \) so that (17) can be written as

\[
k^2 = 2XB - 2j(RB + GX) \hspace{1cm} (18)
\]

and its right hand side is just below the real axis. Therefore, we have a (lossy) propagation for \( XB > 0 \), i.e., in the frequency range of interest (and also above it, up to a frequency of several tens of MHz, where \( X \) becomes negative).

For \( \Delta z \to 0 \), Eq. (13) becomes the standard solution of telegrapher equations [42], with the propagation constant given by (18) and characteristic impedance equal to

\[
Z_0 = \frac{1}{y_1} \simeq \frac{B}{-jk\Delta z} \simeq \frac{2R + jX}{-jk} \\
\simeq -\frac{2X}{k} + 2j\frac{R}{k} \hspace{1cm} (19)
\]

which could be approximated as \(-\frac{2X}{k}\), since its small imaginary part has no physical effects [43].

**III. PROPAGATION BEHAVIOR IN CC SYSTEMS**

The analytical description of the body channel as a transmission line provides the knowledge of its propagative behavior and characteristic impedances. It is worth noting that this model can describe the propagation of the signal in a given body segment, providing a quantitative basis for the improvement of any BCC device.

**A. LIVE WIRE DEVICE**

As a use-case for the assessment of our BCC model, we chose to consider the Live Wire [35], [44]. Live Wire has been implemented by using discrete components, a choice shared with similar BCC devices [2], [45], [46], rather than developing a custom complementary metal oxide semiconductor (CMOS) integrated circuit [33], [47]–[50]. The Live Wire system is shown in the block diagram of Fig. 4. The details about the design and functioning about the Live Wire device can be found in [35], [44]. In the following we report the essential information required to understand the validation of the circuital representation in Fig. 5 and the experimental communication tests.

The TX is a battery powered module, featuring the STM32L486 microcontroller (ST Microelectronics, Geneve, CH). The TX electrode is connected to a low-impedance (50 \( \Omega \)) general purpose input/output (GPIO) pin of the microcontroller. The Live Wire system adopts a self Synchronized-On-Off-Keying (S-OOK) strategy because of its low average power consumption of by the TX (100 mW on average) [31]–[6], [35], [44], [51]. The TX signal, \( V_{TX} \), is a 670 kHz square wave, scaled by non coherent impulses [44]. The carrier frequency was selected considering the trade-off between the possibility to adopt the system as an aid for blind subjects, while respecting safety regulations and preventing, as much as possible for a CC device, the environmental interferences (e.g., to other electronic devices) [3]. The signal is coupled to the human body by means of a 5.5 cm \( \times 3.6 \) cm copper plate acting as a TX electrode (Fig. 4). Even though the electrode area is quite large (about 20 cm\(^2\)), compared to the active area of commercial gelled electrodes for biopotentials recordings, dry electrodes (e.g., the textile ones)
present areas up to 24 cm² [52] and more [53]. The size is comparable to the whole-area (including the support) of several commercial electrodes too for recording (e.g., COVIDIEN H49P or FIAB F9079), and smaller than several electrodes for functional stimulation. The TX electrode can be in direct contact with the body or garments, and it is short-circuited to the external ground to ensure both the direct and return paths [35], [44]. Downstream, the signal propagates in the human body to the receiver unit thanks to a pick-up electrode made up of a 2 cm × 2 cm copper plate. The size of the TX and RX electrodes were chosen considering the trade-off between the dimension of the ground electrodes and the signal strength, as underlined in [3] (see Sect. SM3 of the Supp. Material). The demodulation and synchronization are performed through a software routine running on the microcontroller, which verifies the I/O V status, thus reducing the overall hardware complexity. The Live Wire system requires one 150 mAh Li-Po battery (nominally 3.6 V) and two ADP1712 (Analog Devices, Norwood, MA, USA) linear voltage regulators for the power supply module.

B. CIRCUITAL REPRESENTATION

The application of the model to the Live Wire device is shown in Fig. 5, where the transmitter voltage (V_{TX}), the receiver impedance (Z_{RX}), composed of a capacitance (C_{RX}) with a resistance (R_{RX}) in parallel and the contact impedances between the electrode and the skin in TX and RX (Z_{e1} and Z_{e2} respectively) are shown.

The model has been used to calculate the RX voltage V_{RX}, which is the actual output of the communication using the human body as transmission medium. The voltage at the RX can be derived moving from the calculation of the series impedance between the TX electrode and the receiver impedance Z_{RX}, i.e.,

$$Z_H = Z_{e2} + Z_{RX} \quad (20)$$

where the impedance of the receiver, shown in Fig. 4, is the parallel between a resistance R_{RX} of 1 MΩ and a capacitance C_{RX}, due to a plastic case which separates the metal electrode and the printed circuit board (PCB). In mathematical terms C_{RX} can be estimated as

$$C_{RX} = \frac{\varepsilon_0 \varepsilon_r A}{t} \quad (21)$$

where the dielectric permittivity of the plastic material is \(\varepsilon_r = 2\), A is the plates area (m²) and t, in m, is the distance between the electrode and the PCB. For our case, the electrode size is 5.5 cm \times 3.6 cm and the plates are 2 mm distant, hence, the capacitance is about 17.5 pF.

Knowing the body impedance \(Z_0\) and the propagation constant \(k\) derived from the periodic TML model, by using (19) and (18), with \(\Delta z = 2\) cm, respectively, the impedance seen from the line, \(Z_{IN}\) can be evaluated as

$$Z_{IN} = Z_0 \frac{Z_H \cos (kd) + jZ_0 \sin (kd)}{Z_0 \cos (kd) + jZ_H \sin (kd)} \quad (22)$$

where \(d\) is the distance between electrodes. Then, the current flowing in the Thevenin equivalent is derived as follows

$$I_0 = \frac{V_G}{Z_{TX} + Z_{e1} + Z_{IN}} \quad (23)$$

where \(Z_{TX}\) is the transmitter impedance, defined as

$$Z_{TX} = R_{TX} + \frac{1}{j \omega C_{TX}} \quad (24)$$

The term \(R_{TX}\) is equal to 50 Ω, as experimentally verified, and the radiation losses of the electrode are negligible at the working frequency of 670 kHz. This contribution should be considered due to the ambiguous electrode-antenna behavior observed elsewhere in the literature [3].

The value of the parasitic impedance for the extrinsic return path, \(C_{TX}\), in pF, is a largely variable quantity (from tenths pF to tens of pF), according to the literature, and, furthermore, it has been often considered as a function of the working frequency and of the size of the TX or RX electrodes [20], [31]. However, its explicit dependence from the environment and the working configurations is still an open issue. Provided that, for CC systems, the environments has a considerable influence on the secondary path, it is reasonable to assume that the parasitic capacitance depends on the TX-RX distance \(d\), in m, employing the following second-order polynomial function

$$C_{TX}(d) = c_0 + c_1 d + c_2 d^2 \quad (25)$$

where \(c_2\) is equal to 4 pF/m³, \(c_1\) is 3 pF/m² and, finally, \(c_0 = 1 \text{ pF/m}^{-1}\). The assumption of \(C_{TX} = C_{TX}(d)\) should be interpreted considering that as TX-RX distance increases, the effective area available for the capacitive return path increases too [3]. Different models were tested, namely a constant capacitance value, a first, second and third order polynomial function. However, among the tested functions, the rate of variation of \(C_{TX}(d)\) which allows to best fit the experimental data was Eq. (25).

Then, the voltage in the equivalent circuit is derived as

$$V_0 = Z_{IN} I_0 \quad (26)$$

Hence, the voltage at the end of the body segment can be evaluated with the following transport equation

$$V(d) = V_0 \cos (kd) - jZ_0 I_0 \sin (kd) \quad (27)$$

The current flowing from the transmission line is

$$I(d) = I_0 \cos (kd) - \frac{1}{Z_0} V_0 \sin (kd) \quad (28)$$

Finally, the voltage on the receiver unit is

$$V_{RX} = Z_{RX} I(d) \quad (29)$$

The voltage \(V_{RX}\) has been derived from the proposed model by calculating it using the circuital representation in Fig. 5 and relying on the periodic TML representation of the human body. The obtained quantity was compared to the experimental results in order to validate the model. It must be
pointed out that the quantities $Z_{TX}$ and $Z_{RX}$ are known or can be estimated, whilst the value of $Z_{e1}$ and $Z_{e2}$ depend on the type of electrodes used and on the subject skin. Furthermore, the value of $Z_0$ cannot be estimated from the model unless the line parameters, i.e., the single cell RLC components, are known or found from experimental analysis.

IV. EXPERIMENTALS

A. IMPEDANCE MEASUREMENTS

In order to validate the TLM and circuit representation proposed for the signal transmission, the possible role on the propagation of the impedances of both the electrodes and the dielectric properties of the biological medium were investigated.

![FIGURE 6. Experimental setup for the measurement of the forearm impedance.](image)

To cover the range of working frequencies of BCC, the contact bio-impedance measurements were performed from 50 kHz to 1 MHz, with a 50 kHz span, using a HP 4192A (Keysight Technologies, Santa Rosa, CA, USA) impedance meter, as shown in Fig. 6. The instrument was calibrated at the upper frequency value of 1 MHz. Two FR4-copper electrodes were placed on the anterior forearm, with the metal faces in contact with the skin. The TX and RX are not present in this configuration, thus implying that the measured impedance is the effective impedance without considering $V_G$ (voltage generator), $Z_{e1}$, $Z_{e2}$ in Fig. 5. In particular, one electrode is placed near the wrist and the other one near the elbow, as shown in Fig. 6. On ten voluntary subjects, the impedance was computed as the mean over three consecutive measurements, to smooth over noise and measurement errors. The anthropometric characteristics of the subjects tested in bio-impedentiometric measurements are reported in Tab. 1. The study was performed following the principles outlined in the Helsinki Declaration of 1975, as revised in 2000. Subjects were informed about the procedures, risks, and benefits of the study and that they could withdraw at any time.

![TABLE 1. Biometric data of the subjects tested in bio-impedentiometric measurements.](table)

| Sub. ID | Gender | Age (y) | Height (cm) | Weight (kg) |
|---------|--------|---------|-------------|-------------|
| S1      | M      | 25      | 187         | 82          |
| S2      | F      | 24      | 165         | 60          |
| S3      | M      | 25      | 183         | 65          |
| S4      | F      | 18      | 162         | 57          |
| S5      | F      | 25      | 160         | 60          |
| S6      | F      | 24      | 162         | 58          |
| S7      | M      | 24      | 175         | 68          |
| S8      | F      | 24      | 160         | 62          |
| S9      | M      | 25      | 180         | 70          |
| S10     | M      | 25      | 170         | 71          |

B. CHANNEL CHARACTERIZATION AND TRANSMISSION MEASUREMENTS

The BCC is assessed by measuring the Signal-to-Noise Ratio (SNR) and the communication performances in several cases, as shown in Fig. 7 [54]. The model is used to understand and interpret these experimental situations. The Live Wire device is used in the experiments [35], [44].

This measurement requires to know the voltage at the receiver and noise voltages. To make these data available to the user, RX sends these informations via Bluetooth to the PC. The receiver output is a voltage encoded in 12 bits ($V_{12bit}$), so it is converted in volt ($V_{\text{Volts}}$) in post-processing as follows:

$$V_{\text{Volts}} = \frac{V_{12bit}V_{\text{REF}}}{(2^{\text{12bit}} - 1)}$$

where $V_{\text{REF}}$ is voltage reference of the ADC (2.83 V) and the signal amplification is $A = 22$.

SNR in dB is then measured as:

$$\text{SNR} = 10 \log \left( \frac{P_s - P_n}{P_n} \right) = 10 \log \left( \frac{(V_s - V_n)^2}{R_{RX}} \right)$$

where $P_s$ and $P_n$ are signal and noise power respectively, whilst $V_s$ and $V_n$ are signal and noise voltage respectively.
$V_o$ is voltage offset of the RX ADC and $R_{RX}$ is the input resistance of the receiver, equal to 1 MΩ.

The SNR is a simple method to measure the signal level and it allows to derive information about path gain, however it is not strictly representative of the actual communication performances of the devices, the packet error rate (PER) and total error rate (TER) were analyzed as figure of merits for the actual communication between TX and RX. In the specific four different scenarios may occur, i.e. the packet CRC is correct (packet correctly received), incorrect CRC but correct packet length (packet error), only the first or central part of the packet is received (synchronization error), or the packet is not received. In order to assess TER and PER the following relations have been used:

$$\text{PER} = \frac{\# \text{ Uncorrect Packets}}{\# \text{ Correct Packets}}$$

$$\text{TER} = 1 - \frac{\# \text{ Correct Packets}}{\# \text{ Total packets sent}}$$

For PER calculation only packets with incorrect CRC and correct length are considered. Moreover, in these tests the total packets sent were 500.

The following scenarios have been evaluated for the SNR, TER and PER, namely:

- The near field coupling in air was tested by placing the TX and RX on a table at a variable distance $d_5$, as shown in Fig. 7(a);
- The difference between air communication and the case where the subject’s body acts as transmission medium was investigated for the same distance, as depicted in Fig. 7(b);
- The on-body propagation was studied when the RX was worn on the subject’s arm, by varying the distances between TX and RX, as presented in Fig. 7(c);
- The SNR levels and the transmission were assessed when the signal can be shielded or shunted [44]. Four distinct case were considered. At first, it was tested the case when the subjects wear the RX on their chests and touch themselves with their own hand while contacting the TX, as shown in Fig. 7(d) (Self Hand, No Touch). The effects of the body shielding were studied when another subject touches the subject equipped with the RX, as shown in Fig. 7(e) (Self Hand, Touch). Then, it was investigated the configuration in which this second subject covers the RX with his hand, as shown in Fig. 7(f) (No Self Hand, No Touch). Finally, it was assessed the case when the second subject covers the RX with the hand while touching the subject in...
FIGURE 8. Modulus and phase of the impedance between the TX and RX electrodes placed on the anterior forearm at a distance of 2 cm, 5 cm and 10 cm.

contact with TX, as shown in Fig. 7(g) (No Self Hand, Touch). The distance between the RX and the hand was varied up to 15 cm in a direction normal to the metal surface of the RX;
• In Fig. 7(h), it is represented the case of No Touch Propagation, i.e., the possible transmission of signal when one subject is in contact with the TX, but another person, at a distance D, wears the RX. This very critical but relevant case was studied to assess the influence of the capacitive return path varying the distance between the two subjects;
• The attenuation of a chain of people, from one to six, was investigated by employing the configuration shown in Fig. 7(i).

V. RESULTS
A. IMPEDANCE FITTING AND MODEL VALIDATION
The parameters needed to validate the model were derived from the bio-impedances measured as described in Sect. IV-A. The results of the modulus and phase of the impedance measured on ten different subjects for three different inter-electrode distances are reported in Fig. 8. For all distances between TX and RX, the modulus of the impedance between the two electrodes decreases as frequency increases, being almost constant up above 500 kHz, i.e., for the working frequencies of CC systems, such as the Live Wire device. On the other hand, the phase tends to increase monotonically with the frequency. The variability of the complex impedances between subjects is quite noticeable, with a 30% of uncertainty typical of human dielectric properties [34], [55].

From Fig. 8, the unknown model parameters were derived by numerical fitting using the cost function introduced in Eq. (30). The results for the three inter-electrode distances are reported in Tab. 2. They reveal how the spacing between the resistances $R_{e1}$, $R_{e2}$ and the capacitance $C_{e2}$ proportionally increases by increasing the electrode distance. On the other hand, the capacitance $C_{e1}$, $C_b$ and $C_f$ decreased. Furthermore, the variation of $L_b$ can be explained. Indeed, between the values derived for $d = 2$ cm and $d = 5$ cm there is a 0.3% difference, within the tolerance of the GA. The inductance value derived for $d = 10$ cm is about 0.55 times lower than the one retrieved for a body segment a half and a fifth long, thus implying that the induced currents flowing in the skin decreases as the length of the transmission line increases, which agrees with theoretical and full-wave results of other literature works [3], [8], [9], [34]. As regards the values of $R_b$, it must be underlined that this quantity accounts for the ideal electrical current path if the signal does

| Parameter | $d = 2$ cm | $d = 5$ cm | $d = 10$ cm |
|-----------|------------|------------|------------|
| $R_{e1}$, kΩ | 3.14 | 4.65 | 3.93 |
| $C_{e1}$, nF | 2.37 | 1.85 | 0.03 |
| $R_{e2}$, kΩ | 0.32 | 0.68 | 4.11 |
| $C_{e2}$, pF | 37.0 | 68.3 | 2006 |
| $L_b$, nH | 0.975 | 0.978 | 0.599 |
| $C_b$, nF | 2.37 | 1.85 | 0.03 |
| $R_b$, Ω | 9880 | 10 | 532 |
| $C_f$, pF | 40.8 | 98.4 | 38.7 |
FIGURE 9. Values of maximum and minimum voltage recorded at the receiver for different distances between TX and RX. The curves for the min. and max. RX voltages are the average of ten experiments. The distances are considered moving the RX from the arm (minimum of 10 cm) to the chest (maximum 150 cm).

FIGURE 10. The signal propagation in air and through the body analyzed for distances between TX and RX ranging from 0 cm to 50 cm between TX and RX.

not penetrate along the transverse section of the skin [32]. However, even though $R_b$ is supposed to retain this physical meaning, it has been reported that its value is related with the experimental conditions, the type of coupling and the electrodes configuration and hence it has a range of variation of hundreds of ohms [25], [31].

The fitting curves calculated with the results reported in Tab. 2 are shown in Fig. 8. It can be noticed that the average absolute error on the module is $24.38 \, \Omega$, $27.75 \, \Omega$ and $28.74 \, \Omega$ for 2 cm, 5 cm and 10 cm, respectively. The average absolute error on the phase is $3.29 \, \degree$, $2.13 \, \degree$ and $1.88 \, \degree$ for 2 cm, 5 cm and 10 cm respectively. Using the dielectric dispersion model of the dry skin lead to more than $60 \, \degree$ on the phase impedance and about $100 \, \%$ for the modulus. As previously underlined by [3], [32], the hydration effects of the skin causes the phase constant to increase with frequency, as can be observed from Fig. 8, or derived from Eq. (18), causing the degradation of modulated signals, thus limiting the available bandwidth. This is due to the one order of magnitude difference between the electromagnetic properties of the wet and dry skin across the whole frequency range [39], [56]–[59]. Furthermore, it should be stressed that, using the simple RC model of Wang et al. [25], the deviation from the measured $V_{RX}$ is more than $50\%$, and so, if $C_{TX} = \text{const}$, the calculated receiver voltage would decrease with increasing TX-RX distance. As a result, the proposed periodic transmission line model is capable of describing the behavior of the body impedances with respect to frequency and TX-RX distance.

Given the parameters of Tab. 2 and the estimated values of the RX and TX components, the circuit of Fig. 5 was used to calculate the values of voltage at the receiver in different configurations, taking into account the space dependency of the parasitic capacitance between TX and RX. The results obtained by the study of the complete model are shown in Fig. 9, where the average minimum and maximum experimental voltages measured for the ten subject on bare skin are reported. It can be noticed that the value of voltage at the RX site estimated using the developed TML are able to predict $V_{RX}$ which falls within the $\pm 5 \, \text{mV}$ experimental deviations of the measured data. A slight exception is noticed for the case of an inter-electrode distance of $d = 40 \, \text{cm}$, which exceeds the maximum experimental voltage value of only $0.45\%$. This error is probably due to the polynomial approximation of the ground return capacitance of the transmitter (see Eq. (25)). It is important to underline that, to the best of the authors’ knowledge, the dependence of the return path on the distance has been used for the first time in this work and can successfully explain the lack of a direct dependence between RX voltage and distance. The fitted parameters of TL used in this model are fixed and correspond to a TX-RX distance of 10 cm, i.e., five times the size $\Delta z$ of the unit cell. For increasing distances, it would be expected that these parameters change, but the model proposed is robust with respect to these variations. Indeed, if the TLM parameters remain in the expected range of values, the voltage output variations are negligible. Therefore the model developed can reliably describe the BCC communication.

B. SNR AND SIGNAL TRANSMISSION

In the first scenario, the BCC propagation was tested with and without the body, as shown in Fig 7.(a) and 7.(b). The goal of this test was to demonstrate the range of distances enabling capacitive transmission without the presence of the human body, i.e., when the signal propagates via the parasitic capacitance $C_{TX}$ between the RX and the TX, and when the body acts as transmission medium. The comparison of the SNR in the case of Fig. 7(a)-(c) demonstrates that, for small distances below $5 \, \text{cm}$, the signal propagates from the TX to the RX regardless the communication channel (whether it is the air or the body). This is coherent with the physical mechanism of the capacitive coupling and the proposed model. As shown in Fig. 10, in the presence of the body
channel between the TX and RX, the electromagnetic signal can propagate with lower attenuation than in the air channel, thus realizing a guided-like propagation. When the distance increases, in the case of on-air propagation, the signal levels quickly decrease to zero. Indeed, the transmission does not occur at all for \( d > 5 \) cm. On the other hand, the on-body communication retains the signal attenuation around 35 dB for all the distances analyzed in this test. These results can be explained by using the proposed model and the response of the circuit scheme shown in Fig. 5. If the periodic transmission line representing the body is removed from the circuit in Fig. 5, the signal in the air path attenuates as distance increase (5 dB/cm) due to the increasing value of the parasitic capacitance \( C_{TX} \), according to Eq. (25). This results in a drastic reduction of the SNR. The presence of the body establishes an almost constant signal levels at the RX site (about 40 dB ± 0.12 dB).

Remembering that the BCC device under analysis was developed as a tool for blind subjects [35], it is necessary to investigate the influence of the environment and the possible disturbing effect of body configurations or the interaction with other subjects. Therefore, the effect of body shielding in different configurations was evaluated for the cases presented in Fig. 7(d)-(g). The signal attenuation, in terms of SNR, was assessed in four cases when the same subject or another one place their hand over the RX. The results are shown in Fig. 11. The SNR remains almost constant (in the range between 37-40 dB) in the case the hand touching RX belongs to another person. On the other hand, when the person touching TX lays the own hand over the RX electrode, a shielding effect appears, causing up to 22 dB of reduction in the SNR. When the distance between the RX and the covering hand increases, the SNR levels returns to 40 dB. The SNR increases by 100% in a 10 cm range when the hand of the subject wearing the RX is moved away from the TX. According to our transmission line model, the 50% reduction in SNR values observed in the cases of self-touch can be explained considering that the presence of the hand on the RX induces an anti-phase signal whose propagation degrades the voltage wave. The influence of the destructive signal decreases as the TX-distance increases. Furthermore, when the second person touches the subject a leakage path is created and, as a consequence, part of the signal is de-phased and attenuated. Without taking into account the propagative nature of the body channel, i.e., using simple RC models as found often in the literature [31], the explanation of these cases would be non-trivial.

Given the advantages of BCC, since the practical use of any CC device, such as the case-use Live Wire, is the final application, a series of critical cases were investigated. In particular, the characterization of BCC devices operating in presence of other subjects is an aspect poorly analyzed in the literature [3]–[6]. Therefore, in this work, we investigated the case of No Touch Propagation, i.e., when a subject touches the TX, but another person at a given distance wears the RX, without being in contact, as shown in Fig. 7(h). The SNR values measured moving the subject with the RX in the four directions around the TX, for distances from 5 cm to 120 cm, are presented in Fig. 12. In all the four direction the signal strength reduces while the TX-RX distance increases. The SNR values are not defined for \( d > 90 \) cm, and hence no communication occurs. From Fig. 12 it can be noticed that the SNR varies of \( \sim \pm 12 \) dB in the four directions, which implies that the environment and the return path through the parasitic capacitance may have a strong influence on the body channel propagation.

The results for the chain from one to six people are shown in Fig. 13. The contact between different people does not significantly degrade the signal. In fact, it retains an SNR between 30 dB and 40 dB, which is sufficient to discriminate the signal from the noise at the RX site. As a matter of
fact, the contact between different people degrades the signal. Considering our periodic transmission line model (see Fig. 3), in the case of the people chain, the resulting propagation channel is composed by unit cells having a slightly variable skin properties, which determine this reduction in the signal strength.

C. COMMUNICATION PERFORMANCES

The experimental assessment of the SNR allows to measure the signal levels, given a constant noise level. The results in Fig.s 10-13 provided a useful feedback about the capability of the proposed mathematical framework for describing the physical mechanism of BCC. However, the SNR does not provide information about the communication performances. Therefore, in order to assess the actual communication between TX and RX, the PER and TER are evaluated in the same experimental configurations (see Fig. 7).

The quantification of PER and TER for the case of air vs. body transmission (i.e., cases of Fig. 7(a)-(c)) confirms the necessity of the body presence to ensure the communication, coherently with the proposed model, as can be inferred from Fig. 14. Indeed, in a way similar to the SNR findings from Fig. 10, when the transmitter and the receiver are in contact to each other (e.g. less than 1 cm of distance) the signal can couples by the electrostatic near-field in air and through human tissues. However, when the distance slightly increases and overcome 5 cm, the communication is guaranteed by the guided-transmission in the body. Therefore, the PER in the air channel can be only evaluated when TX and RX are in contact as shown in Fig. 14(a). When the signal is confined in the skin layer, due to the relatively constant SNR levels, for distances up to 15 cm, the PER reaches a maximum of 0.2. As regards the TER values, shown in Fig. 14(b), it can be noticed that the TER is maximum for TX-hand distance up to 2 cm, and then reduces of more than 60%. This implies that the number of correct packets, with respect to the total number of packet sent, is lower than that achieved in the No Self Hand cases, i.e. Fig. 7(h) and 7(i). For the cases of body shielding, the TER values (Fig. 15(b)) show an overall downward trend with increasing distance. It emerges that the hand positioning on the receiver has an effect in all the configurations analyzed, but when the hand-RX distance is zero, or very small, such effect still enables communication only if the hand touching RX is of a second person and no touch occurs. When the hand is distanced more than 5 cm TER values are in the expected range below 40%.

In the most critical cases of No Touch Propagation tests, i.e., the cases Fig. 7(e)-(g), interesting differences across the configurations arises. In particular, the PER values are comprised between 0 and 0.2, as shown in Fig. 15(a). The highest value of 0.6 is reached when the shielding is due to the self hand of the subject at a distance of 2 cm from the TX. Despite this outlier, the PER attenuates with the TX-hand distance (about a 50% reduction in 3 cm). As regards the TER values, shown in Fig. 15(b), it can be noticed that the TER is maximum for TX-hand distance up to 2 cm, and then reduces of more than 60%. This implies that the number of correct packets, with respect to the total number of packet sent, is lower than that achieved in the No Self Hand cases, i.e. Fig. 7(h) and 7(i). For the cases of body shielding, the TER values (Fig. 15(b)) show an overall downward trend with increasing distance. It emerges that the hand positioning on the receiver has an effect in all the configurations analyzed, but when the hand-RX distance is zero, or very small, such effect still enables communication only if the hand touching RX is of a second person and no touch occurs. When the hand is distanced more than 5 cm TER values are in the expected range below 40%.

Finally, the evaluation of PER and TER was performed for the case of people chain. From the findings of Fig. 16, it can be underlined that the PER oscillates around 0 and 0.1 for an increasing number of people and hence, the number

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**FIGURE 13.** SNR levels measured for a chain of one to six people, in contact, interposed between the TX and RX.

**FIGURE 14.** a) Comparison of the PER for the air and BCC transmission mechanisms vs. the TX-RX distance. b) Comparison of the TER for the air and BCC transmission mechanisms vs. the TX-RX distance.
of uncorrect packets about a tenth of the number of correct ones. However, as result from the TER values, the number of correct packets over the total number of packet sent, decreases as the number of people increases. All these findings underline that it is possible to establish the communication using the body as transmission medium, and the case-use device *Live Wire* is capable of that, but still exist some critical cases which may hamper the use of BCC in practice. These limitations can be understood and studied using the proposed mathematical framework to design or enhance the performances of a given CC device for BCC.

**VI. CONCLUSION**

This work dealt with the problem of the development and validation of a transmission line model for BCC. A novel periodic transmission line model was developed in order to describe the signal propagation and to establish a quantitative and general basis for BCC. The unknown model parameters were estimated by fitting the experimental data of bio-impedances measurements performed on ten subjects. The novel finite periodic TLM allowed to describe, from 50 kHz to 1 MHz, the frequency response of the body impedances with average errors of 4° and 30 Ω for the phase angle and impedance modulus, respectively. The model was successfully evaluated by testing the case-use capacitive ground-referred device called *Live Wire* finding good correlation between experimental and calculated data. The agreement is within ±3%, which is remarkable, given the intrinsic variability and the 30% uncertainty associated to the parameters related to signal transmission. The proposed model could be used to enhance the performances of capacitive BCC devices. The signal propagation through the human body channel was characterized using the *Live Wire* device in several scenarios. The communication performances were quantified in terms of SNR. The differences between the over-the-air channel transmission and the on-body propagation were analyzed. The BCC mechanism ensure a 40 dB SNR level when the body channel acts as a wire conveying the information from the TX to the RX. Furthermore, the signal decay was tested when a chain of people constitutes the channel, finding a 12.5% reduction per three people. Some critical case are considered. In particular, the body shielding (or shunt) effects and the No Touch Propagation are commented and discussed relying on the proposed model.

The proposed TLM model is a novel distributed model which account for the finite size of the body channel while including the hardware features of TX and RX. Several key points of this paper need to be further addressed in future works. The periodic transmission line model of the skin and the related circuital representation could be used to design impedance-matching networks at both the TX output and the RX input [3], to enhance the overall performance of a BCC system. As underlined by both the model and the experimental analysis, the proposed mathematical framework can be used to investigate the critical cases of communication between TX and RX in CC systems. As an example, the proposed model can be used to investigate frequency hopping techniques for the performance enhancement [6], [33], [60]. Moreover, since BCC is intended for wearable biomedical applications the type electrodes could be modified to turn the device to a truly wearable piece of electronics. An interesting and
appealing opportunity is represented by the investigation of textile electrodes, either metallic or polymer based [61], [62], already used for biopotentials acquisition [55], [63], [64]. In particular, the polymer-based electrodes are interesting because of their non-metallic nature which could hamper the signal transmission at the selected working frequencies. Finally, the proposed circuitual scheme and the TLM model can be used to investigate this scenario.

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M. B. Lodi et al.: Periodic Transmission Line Model for Body Channel Communication
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