Self-organised quantum dots in marginally twisted MoSe$_2$/WSe$_2$ and MoS$_2$/WS$_2$ bilayers

V. V. Enaldiev,$^{1,2}$ F. Ferreira,$^{1,2}$ J. G. McHugh,$^{1,2}$ and V. I. Fal’ko$^{1,2,3}$

$^1$University of Manchester, School of Physics and Astronomy, Manchester M13 9PL, United Kingdom
$^2$National Graphene Institute, University of Manchester, Manchester M13 9PL, United Kingdom
$^3$Henry Royce Institute, University of Manchester, Manchester M13 9PL, United Kingdom

Moiré superlattices in twistronic heterostructures are a powerful tool for materials engineering. In marginally twisted (small misalignment angle, $\theta$) bilayers of nearly lattice-matched two-dimensional (2D) crystals moiré patterns take the form of domains of commensurate stacking, separated by a network of domain walls (NoDW) with strain hot spots at the NoDW nodes. Here, we show that, for type-II transition metal dichalcogenide bilayers MoX$_2$/WX$_2$ (X=S, Se), the hydrostatic strain component in these hot spots creates quantum dots for electrons and holes. We investigate the electron/hole states bound by such objects, discussing their manifestations via the intralayer intraband infrared transitions. The electron/hole confinement, which is strongest for $\theta < 0.5^\circ$, leads to a red-shift of their recombination line producing single photon emitters (SPE) broadly tuneable around 1 eV by misalignment angle. These self-organised dots can form in bilayers with both aligned and inverted MoX$_2$ and WX$_2$ unit cells, emitting photons with different polarizations. We also find that the hot spots of strain reduce the intralayer MoX$_2$ A-exciton energy, enabling selective population of the quantum dot states.

The formation of minibands is a common moiré superlattice (mSL) effect [1–14], often related to a rigid rotation of one 2D crystal against the other. In general, the approximation of a rigid interlayer twist is valid for lattice-mismatched crystals or larger twist angles, where a short mSL period prohibits the formation of energetically preferential stacking domains of the two crystals. In contrast, for marginally (small-angle) twisted bilayers of crystals with very close lattice constants, the long period of the mSL offers sufficient space for creating preferential stacking areas. That is, the energy gain due to better adhesion can surmount the cost of intralayer strain in each of the constituent crystals. The reconstruction of small-angle twisted bilayers into a network of domains [15] [2H for antiparallel (AP) and 3R for parallel-oriented (P) bilayers] has been observed in various bilayers of transition metal dichalcogenides (TMDs) [16–20]. The observed [16–20] and theoretically modelled [15, 21, 22] structures feature hexagonal (for AP) and triangular (for P) NoDW with nodes hosting few nanometer areas of "chalcogen-on-chalcogen" stacking ($X_iX_b$), which are hot spots of the intralayer strain.

Below we study the effects produced by these hot spots of strain in marginally twisted same-chalcogen heterobilayers MoX$_2$/WX$_2$. While domains/NoDW form in both homo- and heterobilayers, the in-plane intralayer deformations, $u(r)$, in those two systems are qualitatively different. In homobilayers the formation of preferential stacking domains is brought about by twisting locally the crystals toward each other. As a result, the deformations at the domain walls are predominantly shear in character, that is, with $\text{div } u \equiv u_{ii} \to 0$, where $u_{ij} \equiv \frac{1}{2}(\partial_i u_j + \partial_j u_i)$ is a 2D strain tensor, and $u_{ii}$ is its trace. In perfectly aligned ($\theta = 0^\circ$) heterobilayers, lattice mismatch ($\delta \approx 0.2\%$ for MoS$_2$/WS$_2$ and $\delta \approx 0.4\%$ for MoSe$_2$/WSe$_2$) requires an adjustment of the MoX$_2$ ($u_{ii}^{Mo} \approx -\delta$ compression) and WX$_2$ ($u_{ii}^{W} \approx \delta$ expansion) lattices towards each other inside the large area domains. This inflicts a few percent of hydrostatic compression of WX$_2$ and expansion of MoX$_2$ in $X_iX_b$ areas (NoDW nodes), quantified in Figs. 1 and 2.

Here, we single out the hydrostatic strain component because of the critical role it plays in determining the K-valley energies in MoX$_2$/WX$_2$ crystals. Several earlier experimental and density functional theory (DFT) studies [23–25] have agreed that conduction and valence
band edges in TMD monolayers are strongly shifted by hydrostatic strain, but without much sensitivity to shear deformations. This trend is illustrated in Fig. 1. The corresponding shifts of conduction/valence band edges in MoX₂/WX₂ determine the energy of the interlayer exciton (IX). The sign reversal of strain produces a much stronger effect on the band energies in XₐXₕ nodes as compared to the bulk of the domains, Figs. 1 and 2, setting up conditions for electron hole confinement and, consequently, for the appearance of SPE at the NoDW nodes.

The intralayer strain (uᵢʰ⁻Ⅿ) maps in Fig. 2 were computed using a multiscale modelling approach [15], tested in the detailed comparison with the STEM microscopy data [16]. This approach starts with the computation of stacking-dependent MoX₂/WX₂ adhesion energy, W, followed by the parametrization of interpolation formulae [15] for its dependence on the interlayer lateral offset, r₀. For both P and AP bilayers, energetically favourable stackings are those with the largest lateral separation between chalcogens. These stackings are MoXₐ and XₐWXₕ for P-orientation and 2H for AP-orientation. Note that XₐXₕ stacking is unfavourable energetically, and its interlayer distance swells [15] by up to ±0.5 Å. By combining interpolation formulae for W(r₀) [15] where we use local lateral offset,

\[ r₀(r) = \delta \cdot r + \Delta(r), \]  

with elasticity theory and minimizing total energy of the bilayer across its mSL (which period, l ≈ a/√θ² + δ², is fixed by the twist angle and lattice mismatch between the crystals), we compute the deformation fields.

In Fig. 2(a,c), we present uᵢʰ⁻Ⅿ maps for MoSe₂/WSe₂ bilayers with θ = 0° and θ = 0.4°, respectively. The first term in Eq. (2) represents the effect of hydrostatic strain (similar to those in twisted homobilayers). To incorporate strain into the shifts of conduction/valence band edges, δɛᵢ/c/v, we performed DFT modelling of the TMD band structures using Quantum ESPRESSO [26]. For each material, we considered biaxial strain in the range of ±2%, fully relaxing atomic positions in the monolayer with Vanderbilt PBE GBRV ultrasoft pseudopotentials [27], a wavefunction cut-off of 50 nm, and a 20×20×1 k-point grid, sampled according to the Monkhorst-Pack algorithm [28]. Spin-orbit coupling was included by a norm-conserving fully-relativistic psilary PBE PAW pseudopotentials [29] with E_cut = 80 Ry.

The computed variations of all band edges can be described as \( V_{c/v}^{W/Mo} - V_{c/v}^{Mo/W} \), which are quoted for the relevant bilayer bands. Using these values, we compute,

\[ \delta \varepsilon_{c/v}(r) = V_{c/v}^{W/Mo}(r) - e\phi^{W/Mo}(r) \pm \frac{1}{2} \Delta(r), \]  

taking into account strain-dependent piezoelectric potential [15], \(-e\phi^{piezo}(r)\), and offset-dependent potential drop, \( \Delta(r) \), due to interlayer charge transfer (for details see section S1 in Supplementary Information (SI)). The first term in Eq. (2) represents the effect of hydrostatic component of the intralayer strain which was missed in the previous analysis of the same systems [30]. The twist-angle dependences of the computed band edge energies δεᵢ/c for three selected stacking areas (MoXₐ, XₐWXₕ, XₐXₕ for P and 2H, MoₐWy_b, XₐXₕ for AP) are plotted in Fig. 2(b). These figures suggest that XₐXₕ nodes see section S1 in Supplementary Information (SI)).
regions are potential wells for electrons and holes and these wells are the deepest for $\theta \approx \delta$. Based on that we describe the NoDW nodes as trigonally warped quantum dots (QDs), with band edge profiles exemplified in Fig. 3. These QDs are sufficiently deep to accommodate at least two size-quantized states for electrons/holes which retain their distinct $s$ ($L_z = 0$) and $p$ ($L_z = \pm 1$) characteristics due to the $\hat{C}_3$-symmetry of the dots.

Note that QD formation in marginally twisted structures qualitatively differs from band energy profiles in stronger misaligned P-bilayers with $\theta \geq 2^\circ$, where the band edges at K-valley shift into Mo$_X$X$_b$ stacking areas, see Extended data Fig. 1. This crossover agrees with the findings of Refs. [31–33]. In addition, for MoS$_2$/WS$_2$ with $\theta \approx 1.8^\circ$ and MoSe$_2$/WSe$_2$ with $\theta \approx 2.4^\circ$ the energy profile for interlayer interband exciton resembles an antidot superlattice more than an array of QDs. This contrasts with the persistence up to $\theta \approx 3.5^\circ$ of shallow QD arrays for both electrons and holes based at $X_tX_b$ areas in AP-bilayers.

Spectral features of the interlayer interband emissions of self-organised QDs of marginally twisted bilayers are sketched on the bottom insets in Fig. 3, color coded according to the emitted photon polarization, listed in Table 1. Energy separation, $\delta E$, between the QD transition and the iX inside the domains was computed as,

$$\delta E = \varepsilon_e^{(s)} - \varepsilon_h^{(s)} - E_{iX} - \int \int d^2r d^2r' \left| \psi_e^{(s)}(r) \right|^2 \left| \psi_h^{(s)}(r') \right|^2 V_{eh}(r - r').$$

Here, $E_{iX}$ is the iX energy and $\varepsilon_{e/h}^{(s)}$ are the energies of the electron/hole $s$-states $\psi_{e/h}^{(s)}$ inside quantum well. We also take into account the interlayer e-h attraction, $V_{eh}$, screened by the in-plane susceptibility of TMDs and hBN environment [34] (see details in Sections S2, S3 in SI).

The computed dependences of $\Delta E(\theta)$ for MoX$_2$/WX$_2$ bilayers (X=Se,S) are shown on the middle bottom panel in Fig. 3. We find that the QD line can be tuned across a 0.8–1.2 eV spectral interval (telecom range) for $0.3^\circ \leq \theta \leq 1^\circ$ in MoSe$_2$/WSe$_2$ and for $0^\circ \leq \theta \leq 0.5^\circ$ in MoS$_2$/WS$_2$. In Fig. 3 the computed data for electrons in AP-MoS$_2$/WS$_2$ are terminated at $\theta = 0.6^\circ$, because for a larger misalignment the K-valley conduction band profile starts resembling an antidot lattice with maxima at the 2H domains (see Fig. 2(b)). Additional information, displayed in Fig. 3, concerns the fine structure related to a $\Delta_{SO}$-splitting between the spin-flip (SF) and spin-conserving interband transitions inside the QDs, prescribed by spin-orbit splitting in c/v bands of MoX$_2$/WX$_2$ [35–39]. Note that, in each of the two $\pm K$-

Figure 3. Self-organised quantum dots and spectral features of SPE and iX. (Top) Conduction (c) and valence (v) band edge profiles in vicinity of $X_tX_b$ nodes of NoDW in reconstructed P- and AP-MoX$_2$/WX$_2$ bilayers with $\theta = 0^\circ$. Colors of wavy lines encode polarisations of emitted light in $\pm K$-valleys: red for circular and green for $z$-polarisation. Upper/lower subscript of circular polarisation ($\sigma_- or \sigma_+$) indicates helicity of light emitted in $+$K/−K-valleys. Left and right bottom panels show sketches of predicted optical spectra in marginally twisted P- and AP-MoX$_2$/WX$_2$ bilayers, respectively. Middle bottom panel shows the calculated shift of the SPE energy (3) with respect to energy of iX inside of Mo$_X$X$_b$/2H domains for P- and AP-bilayers.
valleys, the hole spin at the band edge is determined by the spin-valley locking in WX₂, whereas conduction band in MoX₂ is characterized by spin-orbit splitting ΔSO. Also, the iX emission from the inner part of domains, shown in Fig. 3, differs for P and AP bilayers. For AP-bilayers we expect a single line of circularly polarised iX emission. For P-bilayers iX energies and polarisations are different for Mo₆X₆ and X₅W₅ domains, with the energy splitting determined by the interlayer charge transfer [30, 40, 41] and circular (in Mo₆X₆) vs linear (in X₅W₅) polarization, established in Ref. [42]. To mention, DFT modelling suggests that lattice matching inside domains promotes direct-to-indirect band gap crossover for electrons towards Q-valley and most importantly for holes towards Γ-valley [43] (see also Section S6 in SI). Also in Table I we highlight the most prominent SPE transition which happens to be related to the spin-conserving recombination in self-organised QDs in P-bilayers.

In addition, intraband s − p optical transitions for electrons/holes trapped in QDs give rise to infrared (IR) features with energies shown in Fig. 4(a) for twist angles 0° ≤ θ ≤ 1°. In AP-bilayers, X₅X₆ NoDW nodes also feature spikes of pseudomagnetic field B∗ [15] characteristic of multivalley semiconductors with a strongly inhomogeneous strain [44–46]. These pseudomagnetic fields have opposite signs for electrons in ±K-valleys, splitting (by \( h\Delta\nu = 5\mu_B B_\star \)) the QD s − p transitions into circularly polarized doublets, as sketched in Fig. 4(a). To mention, such an infrared transition can be used to manipulate the state of the SPE, by exciting either electron or hole into their respective QD p-states.

The intralayer band gap variation, due to the hydrostatic strain at the nodes of domain wall network, generated by the lattice reconstruction in marginally twisted MoX₂/WX₂ bilayers, form a nanoscale array of quantum dots for electrons and holes, which may be operated as single-photon emitters. Based on the presented analysis, we propose that the SPE spectrum can be tuned by the choice of the twist angle over a broad range (including telecom for MoSe₂/WSe₂ bilayers), an the electron/hole state in these QDs can be manipulated via intra-band s − p transitions using THz radiation. The data on the optical oscillator strength of the interlayer interband transitions in such QDs, Table I, suggest that the brightest would be SPEs in marginally twisted biayers with parallel orientation of MoX₂ and WX₂ unit cells. The same data also suggest that the recombination rate of these QD-localised excitons is about ∼ 1% of the recombination rate of the intralayer A-exciton in MoX₂; as the latter was found in Refs. [47–50]

Figure 4. Intralayer infrared features of QDs and intralayer (MoX₂) band gap variation with marking of materials and orientations of the bilayers shown on panel (b). (a) Infrared line, ν, of a QD in P- and AP-MoX₂/WX₂ bilayers, with a line splitting due to pseudomagnetic field, B∗, in AP-structures. Twist angle dependence of QD s − p-transition frequency was analysed based on Eq. 2 and B∗ was estimated for conduction band electrons in MoS₂ using parameters from Ref. [44]. Pseudomagnetic field maps around X₆ nodes of NoDW for θ = 0° and θ = 0.6° are shown in the insets. Pseudomagnetic fields of similar magnitude are expected for holes in MoS₂/WSe₂ bilayers and both electrons and holes for MoSe₂/WSe₂ bilayers. (b) Depth of a potential well, \( U_{AX}^\star \), confining intralayer A-excitons in MoX₂, created by hydrostatic strain around X₆ nodes of NoDW. Inset shows anisotropy of the well profile for P- and AP-MoSe₂/WSe₂ bilayers with θ = 0°.

to be ∼ 200 − 300 fs⁻¹, this would set a 100 MHz possible repetition rate for SPEs in these self-organised QDs. Note that areal density of these SPEs is ∼10¹¹ cm⁻², which is 100 times higher than the density of quantum emitters in patterned TMD monolayers [51, 52] and that a red-shift of the A-exciton in MoX₂, due to the same hot spots of strain, would enable selective population of the QD states for the optical pumping of the self-organised SPEs.

Acknowledgements. This work was supported by
Table I. Polarization of emitted photons by SPE, iX and AX in MoX$_2$, and strength of the emission characterised by interband velocity matrix element, $|v_{CV}^\pm|$ (in $10^6 \times$ cm/s, $v_{CV}^\pm = v_x^\pm \pm iv_y^\pm$) computed with Quantum ESPRESSO (for details see Sections S4 and S5 in SI). $\sigma_+$ ($\sigma_-$) designates photons with clockwise (counter clockwise) polarization outgoing in positive direction of spin quantization axis. $z$ stands for linear out-of-plane polarization. Here, upper (lower) index corresponds to emission from $+K$ ($-K$) valley. For this analysis we employed angular momentum conservation for K-point Bloch states (we used center of coordinate set at a chalcogen atom).

|          | AX MoX$_2$ | P | AP |
|----------|------------|---|----|
|          | SPE$_{X_b}$, iX$_{X_b}$, iX$_{X_b}$ | iX$_{X_b}$, iX$_{X_b}$ | SPE$_{X_b}$, iX$_{X_b}$ |
| ground ($E$) | $\sigma_\pm$ | $\sigma_\pm$ | $\sigma_\pm$ | $\sigma_\pm$ |
| S        | 113        | 6.42 | 4.57 | 2.53 | 0.48 | 2.28 |
| Se       | 97         | 7.03 | 5.30 | 3.73 | 0.54 | 3.49 |
| excited ($E + \Delta_{SO}$) | $\sigma_\pm$ | $\sigma_\pm$ | $\sigma_\pm$ | $\sigma_\pm$ | $\sigma_\pm$ | $\sigma_\pm$ |
| S        | 0.49       | 2.82 | 0.12 | 1.18 | 10.30 |
| Se       | 0.74       | 5.01 | 0.10 | 1.89 | 10.64 |

Extended data Fig. 1. (Top) Twist-angle-dependences of interlayer band gap at K-valleys at high symmetry stacking areas of moiré supercells of P- and AP-MoX$_2$/WX$_2$ bilayers with $1^\circ \leq \theta \leq 3.5^\circ$. For P-bilayers with $\theta \geq 2^\circ$, we find a crossover from an array of QDs located in X$_b$ nodes to an antidot array with shallow minima at X$_b$ areas and high peaks at Mo$_b$ areas, while for AP-bilayers there is no such transition. (Bottom) Real space maps showing crossover (towards larger misalignment) of the interlayer band gap for MoS$_2$/WS$_2$ bilayers with $\theta = 2.5^\circ$: an antidot array for P-bilayers and shallow QD array for AP-bilayers.

---

[1] Y. Cao, V. Fatemi, A. Demir, S. Fang, S. L. Tomarken, J. Y. Luo, J. D. Sanchez-Yamagishi, K. Watanabe, T. Taniguchi, E. Kaxiras, and et al., Correlated insulator behaviour at half-filling in magic-angle graphene superlattices, Nature 556, 80 (2018).
[2] Y. Cao, V. Fatemi, S. Fang, K. Watanabe, T. Taniguchi, E. Kaxiras, and P. Jarillo-Herrero, Unconventional superconductivity in magic-angle graphene superlattices, Nature 556, 43 (2018).
[3] M. Yankowitz, S. Chen, H. Polshyn, Y. Zhang, K. Watanabe, T. Taniguchi, D. Graf, A. F. Young, and C. R. Dean, Tuning superconductivity in twisted bilayer graphene, Science 363, 1059 (2019).
[4] X. Lu, P. Stepanov, W. Yang, M. Xie, M. A. Aamir, I. Das, C. Urgell, K. Watanabe, T. Taniguchi, G. Zhang, et al., Superconductors, orbital magnets and correlated states in magic-angle bilayer graphene, Nature 574, 653 (2019).
[5] S. Xu, M. M. Al Ezzi, N. Balakrishnan, A. Garcia-Ruiz, B. Tsim, C. Mullan, J. Barrier, N. Xin, B. A. Piot,
T. Taniguchi, K. Watanabe, A. Carvalho, A. Mishchenko, A. K. Geim, V. I. Fal’ko, S. Adam, A. H. C. Neto, K. S. Novoselov, and Y. Shi, Tunable van hove singularities and correlated states in twisted monolayer–bilayer graphene, Nature Physics 17, 619 (2021).

[6] A. L. Sharpe, E. J. Fox, A. W. Barnard, J. Finney, K. Watanabe, T. Taniguchi, M. Kastner, and D. Goldhaber-Gordon, Emergent ferromagnetism near three-quarters filling in twisted bilayer graphene, Science 365, 605 (2019).

[7] H. Polshyn, J. Zhu, M. A. Kumar, Y. Zhang, F. Yang, C. L. Tschirhart, M. Serlin, K. Watanabe, T. Taniguchi, A. H. MacDonald, et al., Electrical switching of magnetic order in an orbital chern insulator, Nature 588, 66 (2020).

[8] G. Chen, L. Jiang, S. Wu, B. Lyu, H. Li, B. L. Chittari, K. Watanabe, T. Taniguchi, Z. Shi, J. Jung, Y. Zhang, and F. Wang, Evidence of a gate-tunable mott insulator in a trilayer graphene moiré superlattice, Nature Physics 15, 237 (2019).

[9] S. Chen, M. He, Y.-H. Zhang, V. Hsieh, Z. Fei, K. Watanabe, T. Taniguchi, D. H. Cobden, X. Xu, C. R. Dean, et al., Electrically tunable correlated and topological states in twisted monolayer–bilayer graphene, Nature Physics 17, 374 (2021).

[10] C. Shen, Y. Chu, Q. Wu, N. Li, S. Wang, Y. Zhao, J. Tang, J. Liu, J. Tian, K. Watanabe, et al., Correlated states in twisted double bilayer graphene, Nature Physics 16, 520 (2020).

[11] J. M. Park, Y. Cao, K. Watanabe, T. Taniguchi, and P. Jarillo-Herrero, Tunable strongly coupled superconductivity in magic-angle twisted trilayer graphene, Nature 590, 249 (2021).

[12] Y. Cao, D. Rodan-Legrain, J. M. Park, N. F. Yuan, K. Watanabe, T. Taniguchi, R. M. Fernandes, L. Fu, and P. Jarillo-Herrero, Nematicity and competing orders in superconducting magic-angle graphene, Science 372, 264 (2021).

[13] X. Liu, Z. Hao, E. Khalaf, J. Y. Lee, Y. Ronen, H. Yoo, D. H. Najaafabadi, K. Watanabe, T. Taniguchi, A. Vishwanath, et al., Tunable spin-polarized correlated states in twisted double bilayer graphene, Nature 583, 221 (2020).

[14] A. C. Gadelha, D. A. A. Ohlberg, C. Rabelo, E. G. S. Neto, T. L. Vasconcelos, J. L. Campos, J. S. Lemos, V. Ornelas, D. Miranda, R. Nadas, F. C. Santana, K. Watanabe, T. Taniguchi, B. van Troeye, M. Lamparski, V. Meunier, V.-H. Nguyen, D. Paszko, J.-C. Charlier, L. C. Campos, L. G. Cançado, G. Medeiros-Ribeiro, and A. Jorio, Localization of lattice dynamics in low-angle twisted bilayer graphene, Nature 590, 405 (2021).

[15] V. V. Enaldiev, V. Zolyomi, C. Yelgel, S. J. Magorrian, and V. I. Fal’ko, Stacking domains and dislocation networks in marginally twisted bilayers of transition metal dichalcogenides, Phys. Rev. Lett. 124, 206101 (2020).

[16] A. Weston, Y. Zou, V. Enaldiev, A. Summerfield, N. Clark, V. Zolyomi, A. Graham, C. Yelgel, S. Magorrian, M. Zhou, J. Zultak, D. Hopkins, A. Barinov, T. H. Bointon, A. Kretinin, N. R. Wilson, P. H. Beaton, V. I. Fal’ko, S. J. Haigh, and R. Gorbachev, Atomic reconstruction in twisted bilayers of transition metal dichalcogenides, Nature Nanotechnology 15, 592 (2020).

[17] M. R. Rosenberger, H.-J. Chuang, M. Phillips, V. P. Oleshko, K. M. McCreary, S. V. Sivaram, C. S. Hellberg, and B. T. Jonker, Twist angle-dependent atomic reconstruction and moiré patterns in transition metal dichalcogenide heterostructures, ACS Nano 14, 4550 (2020).

[18] J. Sung, Y. Zhou, G. Scru, V. Zolyomi, T. I. Andersen, H. Yoo, D. S. Wild, A. Y. Joe, R. J. Gelly, H. Heo, S. J. Magorrian, D. Bérubé, A. M. M. Valdivia, T. Taniguchi, K. Watanabe, M. D. Lukin, P. Kim, V. I. Fal’ko, and H. Park, Broken mirror symmetry in excitonic response of reconstructed domains in twisted WS2/MoSe2 bilayers, Nature Nanotechnology 15, 750 (2020).

[19] L. J. McGilly, A. Kerelsky, N. R. Finney, K. Shapovalov, E.-M. Shih, A. Ghiotto, Y. Zeng, S. L. Moore, W. Wu, Y. Bai, K. Watanabe, T. Taniguchi, M. Stengel, L. Zhou, J. Hone, X. Zhu, D. N. Basov, C. Dean, C. E. Dreyer, and A. N. Pasupathy, Visualization of moiré superlattices, Nature Nanotechnology 15, 580 (2020).

[20] S. Shabani, D. Halbertal, W. Wu, M. Chen, S. Liu, J. Hone, W. Yao, D. N. Basov, X. Zhu, and A. N. Pasupathy, Deep moiré potentials in twisted transition metal dichalcogenide bilayers, Nature Physics 17, 720 (2021).

[21] M. H. Naik and M. Jain, Ultraflatbands and shear solitons in moiré patterns of twisted bilayer transition metal dichalcogenides, Phys. Rev. Lett. 121, 266401 (2018).

[22] S. Carr, D. Massatt, S. B. Torrisi, P. Cazeaux, M. Luskin, and E. Kaxiras, Relaxation and domain formation in incommensurate two-dimensional heterostructures, Phys. Rev. B 98, 224102 (2018).

[23] H. J. Conley, B. Wang, J. I. Ziegler, R. F. Haglund, S. T. Pantelides, and K. I. Bolotin, Bandgap engineering of strained monolayer and bilayer MoS2, Nano Letters 13, 3626 (2013).

[24] K. P. Dhakal, S. Roy, H. Jang, X. Chen, W. S. Yun, H. Kim, J. Lee, J. Kim, and J.-H. Ahn, Local strain induced band gap modulation and photoluminescence enhancement of multilayer transition metal dichalcogenides, Chemistry of Materials 29, 5124 (2017).

[25] K. Zollner, P. E. F. Junior, and J. Fabian, Strain-tunable orbital, spin-orbit, and optical properties of monolayer transition-metal dichalcogenides, Physical Review B 100, 195126 (2019).

[26] P. Giannoni, S. Baroni, N. Bonini, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, G. L. Chiarotti, M. Cococcioni, I. Dabo, A. Del Corso, S. de Gironcoli, S. Fabris, G. Fratesi, R. Gebauer, U. Gerstmann, C. Gougoussis, A. Kokalj, M. Lazzeri, L. Martin-Samos, N. Marzari, F. Mauri, R. Mazzarello, S. Paolini, A. Pasquarello, L. Paulatto, C. Sarrao, S. Scandolo, G. Scuﬀer, A. P. Seitsonen, A. Smogunov, P. Umani, and R. M. Wentzcovitch, Quantum espresso: a modular and open-source software project for quantum simulations of materials, Journal of Physics: Condensed Matter 21, 395502 (19pp) (2009).

[27] K. F. Garrity, J. W. Bennett, K. M. Rabe, and D. Vanderbilt, Pseudopotentials for high-throughput DFT calculations, Computational Materials Science 81, 446 (2014).

[28] H. J. Monkhorst and J. D. Pack, Special points for brillouin-zone integrations, Physical Review B 13, 5188 (1976).

[29] We used the relevant relativistic pseudopotentials from http://www.quantum-espresso.org.

[30] V. V. Enaldiev, F. Ferreira, S. J. Magorrian, and V. I. Fal’ko, Piezoelectric networks and ferroelectric domains in twistronic superlattices in WS2/MoS2 and WSe2/MoSe2 bilayers, 2D Materials 8, 025030 (2021).
[31] H. Yu, G.-B. Liu, J. Tang, X. Xu, and W. Yao, Moiré excitons: From programmable quantum emitter arrays to spin-orbit–coupled artificial lattices, Science Advances 3, 10.1126/sciadv.1701696 (2017).

[32] F. Wu, T. Lovorn, and A. H. MacDonald, Theory of optical absorption by interlayer excitons in transition metal dichalcogenide heterobilayers, Phys. Rev. B 97, 035306 (2018).

[33] S. Brem, C. Linderälv, P. Erhart, and E. Malic, Tunable phases of moiré excitons in van der waals heterostructures, Nano Letters 20, 8534 (2020).

[34] M. Danovich, D. A. Ruiz-Tijerina, R. J. Hunt, M. Szyniszewski, N. D. Drummond, and V. I. Fal’ko, Localized interlayer complexes in heterobilayer transition metal dichalcogenides, Phys. Rev. B 97, 195452 (2018).

[35] K. F. Mak, K. He, J. Shan, and T. F. Heinz, Control of valley polarization in monolayer MoS2 by optical helicity, Nature Nanotechnology 7, 494 (2012).

[36] J. P. Echeverry, B. Urbaszek, T. Amand, X. Marie, and I. C. Gerber, Splitting between bright and dark excitons in transition metal dichalcogenide monolayers, Physical Review B 93, 121107 (2016).

[37] G. Wang, C. Robert, M. Glazov, F. Cadiz, E. Courtade, T. Amand, D. Lagarde, T. Taniguchi, K. Watanabe, B. Urbaszek, and X. Marie, In-plane propagation of light in transition metal dichalcogenide monolayers: Optical selection rules, Physical Review Letters 119, 047401 (2017).

[38] G.-B. Liu, W.-Y. Shan, Y. Yao, W. Yao, and D. Xiao, Three-band tight-binding model for monolayers of group-VIB transition metal dichalcogenides, Physical Review B 88, 085433 (2013).

[39] A. Kormányszos, G. Burkard, M. Gmitra, J. Fabian, V. Zdzyomil, N. D. Drummond, and V. Fal’ko, k-p theory for two-dimensional transition metal dichalcogenide semiconductors, 2D Materials 2, 022001 (2015).

[40] F. Ferreira, V. V. Enaldiev, V. I. Fal’ko, and S. J. Magorrian, Weak ferroelectric charge transfer in layer-asymmetric bilayers of 2D semiconductors, Scientific Reports 11, 13422 (2021).

[41] A. Weston, E. G. Castanon, V. Enaldiev, F. Ferreira, S. Bhattacharjee, S. Xu, H. Corte-León, Z. Wu, N. Clark, A. Summerfield, T. Hashimoto, Y. Gao, W. Wang, M. Hamer, H. Read, L. Fumagalli, A. V. Kretinin, S. J. Haigh, O. Kazakova, A. K. Geim, V. I. Fal’ko, and R. Gorbachev, Interfacial ferroelectricity in marginally twisted 2D semiconductors, Nature Nanotechnology 10.1038/s41565-022-01072-w (2022).

[42] H. Yu, G.-B. Liu, and W. Yao, Brightened spin-triplet interlayer excitons and optical selection rules in van der waals heterobilayers, 2D Materials 5, 035021 (2018).

[43] J. Kienle, F. Sigger, M. Lorke, B. Miller, K. Watanabe, T. Taniguchi, A. Holleitner, and U. Wurstbauer, Control of the orbital character of indirect excitons in MoS2/WS2 heterobilayers, Physical Review B 101, 121404 (2020).

[44] H. Rostami, R. Roldán, E. Cappelluti, R. Asgari, and F. Guinea, Theory of strain in single-layer transition metal dichalcogenides, Physical Review B 92, 195402 (2015).

[45] S. V. Iordanskii and A. E. Koshelev, Dislocations and localization effects in multivalley conductors, JETP Lett. 41, 471 (1985).

[46] H. Suzuura and T. Ando, Phonons and electron-phonon scattering in carbon nanotubes, Physical Review B 65, 235412 (2002).

[47] C. Poellmann, P. Steinleitner, U. Leierseder, P. Nagler, G. Plechinger, M. Porer, R. Bratschitsch, C. Schüller, T. Korn, and R. Huber, Resonant internal quantum transitions and femtosecond radiative decay of excitons in monolayer WSe2, Nature Materials 14, 889 (2015).

[48] P. Dey, J. Paul, Z. Wang, C. Stevens, C. Liu, A. Romero, J. Shan, D. Hilton, and D. Karaškaj, Optical coherence in atomic-monolayer transition-metal dichalcogenides limited by electron-phonon interactions, Physical Review Letters 116, 127402 (2016).

[49] T. Jakubczyk, V. Delmonte, M. Kopierski, K. Noga-jewski, C. Faugeras, W. Langbein, M. Potemski, and J. Kasprzak, Radiatively limited dephasing and exciton dynamics in MoS2 monolayers revealed with four-wave mixing microscopy, Nano Letters 16, 5333 (2016).

[50] F. Cadiz, E. Courtade, C. Robert, G. Wang, Y. Shen, H. Cai, T. Taniguchi, K. Watanabe, H. Carrere, D. Lagarde, M. Manca, T. Amand, P. Renucci, S. Tongay, X. Marie, and B. Urbaszek, Excitonic linewidth approaching the homogeneous limit in MoS2 -based van der waals heterostructures, Physical Review X 7, 021026 (2017).

[51] C. Palacios-Berraquero, D. M. Kara, A. R.-P. Montblanch, M. Barbone, P. Latawiec, D. Yoon, A. K. Ott, M. Loncar, A. C. Ferrari, and M. Atatüre, Large-scale quantum-emitter arrays in atomically thin semiconductors, Nature Communications 8, 10.1038/ncomms15093 (2017).

[52] Y. Luo, G. D. Shepard, J. V. Ardelean, D. A. Rhodes, B. Kim, K. Barmak, J. C. Hone, and S. Strauf, Deterministic coupling of site-controlled quantum emitters in monolayer WSe2 to plasmonic nanocavities, Nature Nanotechnology 13, 1137 (2018).