The Fano-Rashba effect

Llorenç Serra¹,² and David Sánchez¹
¹ Departament de Física, Universitat de les Illes Balears, E-07122 Palma de Mallorca, Spain
² Institut Mediterrani d’Estudis Avançats IMEDEA (CSIC-UIB), E-07122 Palma de Mallorca, Spain
E-mail: llorens.serra@uib.es, david.sanchez@uib.es

Abstract. We analyze the linear conductance of a semiconductor quantum wire containing a region where a local Rashba spin-orbit interaction is present. We show that Fano lineshapes appear in the conductance due to the formation of quasi bound states which interfere with the direct transmission along the wire, a mechanism that we term the Fano-Rashba effect. We obtain the numerical solution of the full Schrödinger equation using the quantum-transmitting-boundary method. The theoretical analysis is performed using the coupled-channel model, finding an analytical solution by ansatz. The complete numerical solution of the coupled-channel equations is also discussed, showing the validity of the ansatz approach.

1. Introduction
The spin-orbit (SO) interactions in semiconductor nanostructures are currently attracting much interest due to their central role in the characterization and control of spin properties of these systems [1]. Particular attention is devoted to the Rashba interaction in two-dimensional quantum wells—a SO coupling stemming from the well inversion asymmetry in the growth direction whose intensity can be manipulated by electrical gating [2, 3]. The Rashba interaction is also the underlying physical mechanism for the operation of the proposed spin transistor, one of the paradigms of spintronic device, exploiting the electron’s spin to control the charge current between two ferromagnetic contacts [4].

In a quantum wire with uniform SO interaction the energy subbands deviate from the simple parabolic behavior, with band splittings and anticrossings that yield a modified wire conductance and the possibility of spin-textured states [5]. In this work we consider a wire with nonuniform Rashba interaction, localized to a finite region acting as a spin scattering center (a Rashba dot) that affects the wire conductance properties. It will be shown with numerical and analytical results that this geometry leads in a natural way to the appearance of Fano resonance profiles in the energy dependence of the conductance that depend quite sensitively on the properties of the Rashba dot, such as its dimensions and SO intensity. We refer to this influence of the Rashba SO coupling on the wire conductance as the Fano-Rashba effect.

Fano resonances [6] are a general phenomenon that has been observed in different fields, such as atomic physics [7], Raman scattering [8] and mesoscopic electron transport [9]. Fano-resonance physics appears wherever there is an interference between two paths, one corresponding to direct transmission and the other to the passage through a quasi-bound state lying nearby in energy. As a consequence, characteristic asymmetric lineshapes appear in the most general case, with conductance dips in which transmission is greatly quenched. Indeed, for
the case of scattering centers that may be modelled as attractive potentials it was shown that
an exact transmission zero should always be present [10, 11]. In our case of a Rashba dot the
existence of transmission zeros is not always guaranteed and only for some specific dimensions
and intensities of SO coupling the wire conductance vanishes at the dip position.
A useful theoretical framework to understand Fano resonances is the coupled-channel model
(CCM). In this approach the wave function is approximated by the superposition of different
channel components, each of them obeying an equation containing the channel background
contribution and also the coupling with the other channels. We shall compare the prediction
of CCM with the exact numerical solution obtained from a grid discretization of the Schrödinger
equation. The prediction of the CCM will be analyzed using the ansatz for the quasi-bound
state channel proposed by Nöckel and Stone [11], and also with a full numerical solution of
the CCM equations. In this way we shall explicitly proof the validity of the ansatz. This
analysis is an extension of our recent work [12] and is also related to the study in Ref. [13]
where numerical results for the transmission of wires with a modulated Rashba intensity were
obtained. A localized Rashba interaction in a quantum wire was also considered in Ref. [14] for
the study of shot noise and entanglement within a beam splitter configuration.

2. The physical system
2.1. Hamiltonian
We assume the effective-mass model for the conduction band states of a two-dimensional electron
gas in GaAs and consider a parabolic confinement in the $y$ direction, yielding a ballistic quantum
wire along $x$. The system Hamiltonian reads

$$H = \frac{p_x^2 + p_y^2}{2m} + \frac{1}{2}m\omega_0^2y^2 + H_R.$$  

The oscillator is used to define our energy $\hbar \omega_0$ and length $\ell_0 = \sqrt{\hbar/m\omega_0}$ units. The Rashba
Hamiltonian $H_R$ that we consider in Eq. (1) is characterized by an intensity $\alpha(x)$ vanishing
everywhere except for $0 < x < \ell$ where it takes the value $\alpha_0$. In detail, $H_R$ reads

$$H_R = \frac{\alpha(x)}{2\hbar} (p_y \sigma_x - p_x \sigma_y) + h.c.$$  

where $\sigma_x$ and $\sigma_y$ are the Pauli matrices and the Hermitian conjugation is used to ensure
Hermiticity. The inset to Fig. 1 shows a simple sketch of the physical system under consideration.

2.2. Localized states in ballistic transport
We have obtained the linear conductance $G$ as a function of the electron energy (Fermi energy)
discretizing the $xy$ plane in a uniform grid and using the quantum-transmitting-boundary
algorithm [15]. This method transforms the Schrödinger equation into a complex linear system
whose dimension is twice the number of grid points, due to the two spin components. In spite
of the typically large dimension, the system is highly sparse and it can be solved efficiently with
standard numerical routines. This approach provides the exact wave function for the ballistic
transport problem from which one obtains the transmission amplitudes and the wire’s linear
conductance.

Figure 1 shows a typical conductance curve for specific values of $\alpha$ and $\ell$. The usual staircase
conductance is modified because of the Rashba dot in two main aspects. First the initial part
of the step is smoothed, a usual quantum behavior of scattering through potential wells. The

1 In the numerical applications we take $\alpha(x) = \alpha_0[f(x - \ell) - f(x)]$, with $f(x) = 1/(1 + e^{x/\sigma})$ and $\sigma$ a small
diffusivity. The results are not very sensitive to the precise value of $\sigma$ provided $\sigma < \ell$. 
second conspicuous feature is the existence of a pronounced conductance dip near the end of the conductance plateau. It was already mentioned that the existence of conductance dips is normally due to quasibound states. To proof it in the present context we show in Fig. 2 the wave function for two different energies, \( E_a \) and \( E_b \), where \( E_a \) is chosen to lie in the smooth conductance region while \( E_b \) corresponds precisely to the dip position. Indeed, at energy \( E_b \) the wave function strongly resonates within the Rashba dot, being so amplified that it greatly exceeds the values in the leads. This explicitly shows the dramatic influence of the quasibound state. A systematic analysis of the dip positions for varying \( \alpha \)'s and \( \ell \)'s was performed in Ref. [12].

3. The coupled channel model
3.1. Channel equations
The physics behind the results shown in the preceding section is more easily analyzed within the CCM. Assume the following expansion of the electron wave function \( \Psi(x, y, \eta) \), with \( \eta = \uparrow, \downarrow \), in the transverse oscillator modes \( \phi_n(y) \), with energies \( \varepsilon_n \), and spin eigenstates \( \chi_{s}(\eta) \) of \( \sigma_y \),

\[
\Psi(x, y, \eta) = \sum_{n,s=\pm} \psi_{ns}(x)\phi_n(y)\chi_s(\eta) .
\] (3)

Substituting Eq. (3) into the Schrödinger equation \( H\Psi = E\Psi \) and projecting onto the functions \( \phi_n\chi_s \) we obtain the CCM equations for the different channel amplitudes \( \psi_{ns}(x) \). We shall restrict to the first conductance plateau, i.e., to electron energies fulfilling \( \varepsilon_1 < E < \varepsilon_2 \) and also truncate the expansion in Eq. (3) to the first two transverse modes \( n = 1, 2 \). The selection rules stemming from the Rashba interaction require spin flip and also \( \Delta n = \pm 1 \). Therefore we obtain two independent two-band models, with \( \psi_{1+} \) coupled to \( \psi_{2-} \) and \( \psi_{1-} \) to \( \psi_{2+} \). Actually, the two models are equivalent and we shall focus for simplicity on the \( \psi_{1+}-\psi_{2-} \) system.

The CCM equations for the \( \psi_{1+}, \psi_{2-} \) amplitudes are greatly simplified by means of the following gauge transformation \( \psi_{1+},2- \rightarrow \psi_{1+},2- \exp(\pm i \int k_R(x')dx') \), where we have defined
The exact solution of the preceding section (solid line). As shown by the dash-dotted line of Fig. 1, the ansatz solution satisfactorily expresses in Eq. (7) corresponds to the generalized Fano lineshape, with $$\epsilon \equiv \frac{E - \varepsilon_2 - \Delta}{\Gamma}$$ and $$q = \delta/\Gamma + i(\gamma/\Gamma - 1)$$. In general, the Rashba dot yields complex $$q$$'s and therefore the factor $$|\epsilon + q|$$ in Eq. (7) can not exactly vanish at any $$E$$, thus hindering the formation of zero transmission dips. For some particular values of $$\alpha$$ and $$\ell$$, however, we showed in Ref. [12] that there are accidental cases where the conductance dip is compatible with a zero value within numerical precision. As shown by the dash-dotted line of Fig. 1, the ansatz solution satisfactorily reproduces the conductance dip although, of course, there are some minor discrepancies with the exact solution of the preceding section (solid line).
3.3. Numerical solution

The exact solution of the CCM equations can also be obtained numerically, avoiding the ansatz, using a 1D formulation of the quantum-transmitting-boundary algorithm. Dashed line in Fig. 1 corresponds to the conductance obtained from this numerical CCM approach while Fig. 3 shows the different channel amplitudes (including also the \( \psi_{3+} \) channel), as well as the bound state \( \phi_0 \), for the dip energy \( E_b \) of Fig. 1. The probability density for channel \( \psi_{1+} \) again proves the propagating character of this channel although transmission is rather low at this energy. There is good qualitative agreement of the probability densities \( \rho_2 \) and \( \rho_0 \), explicitly proving that the ansatz is indeed a reasonable assumption for the evanescent channels. The numerical solution has been obtained including also channel \( \psi_{3+} \), however, the smallness of \( \rho_{3+} \) with respect to the other channel densities supports the truncation to the lowest two modes of the preceding subsection.

4. Conclusions

We have theoretically investigated the influence of a localized Rashba scattering center (a Rashba dot) on the linear conductance of a quantum wire. The Rashba dot sustains quasi-bound states that interfere with the direct transmission along the wire and lead to Fano-resonance profiles in the energy dependence of the linear conductance. This Fano-Rashba effect has been elucidated with numerical calculations using a grid discretization of the 2D Schrödinger equation. We also analyzed a coupled-channel model, finding analytical expressions within the ansatz approach that clearly show the appearance of the Fano function for the transmission. Finally, the validity of the ansatz solution and the truncation to the lowest two modes has been assessed by finding the numerical solution of the multimode coupled-channel-model equations.

Acknowledgments

We acknowledge R. López for valuable discussions. This work was supported by the Grant No. FIS2005-02796 (MEC) and the “Ramón y Cajal” program.

References

[1] Zutic I, Fabian J and Das Sarma S 2004 Rev. Mod. Phys. 76 323
[2] Rashba E I 1960, Fiz. Tverd. Tela (Leningrad) 2 1224 [Sov. Phys. Solid State 2 1109]
[3] Nitta J, Akazaki T, Takayanagi H and Enoki T 1997 Phys. Rev. Lett. 78 1335; Engels G, Lange J, Schäpers Th and Lüth H 1997 Phys. Rev. B 55 R1958
[4] Datta S and Das B 1990 Appl. Phys. Lett. 56 665
[5] Schäpers Th, Knobbe J and Guzenko V A 2004 Phys. Rev. B 69 235323; Mireles F and Kirczenow G 2001 Phys. Rev. B 64 024426; Governale M and Zülicke U 2002 Phys. Rev. B 66 073311; Serra Ll, Sánchez D and López R 2005 Phys. Rev. B 72 235309
[6] Fano U 1961 Phys. Rev. 124 1866
[7] Adair R K, Bokelman C K and Peterson R E 1949 Phys. Rev. 76 308
[8] Cerdeira F, Fjeldly T A and Cardona M 1973 Phys. Rev. B 8 4734
[9] Gores J et al 2000 Phys. Rev. B 62 2188; Kobayashi K, Aikawa H, Katsumoto S and Iye Y 2002 Phys. Rev. Lett. 88 256806
[10] Gurvitz S A and Levinson Y B 1993 Phys. Rev. B 47 10578
[11] Nöckel J U and Stone A D 1994 Phys. Rev. B 50 17415
[12] Sánchez D and Serra LI 2006 preprint
[13] Zhang L, Brusheim P and Xu H Q 2005 Phys. Rev. B 72 045347
[14] Egues J C, Burkard G and Loss D 2002 Phys. Rev. Lett. 89 176401
[15] Lent C S and Kirkner D J 1990 J. Appl. Phys. 67 6353