Mechanically Switchable Wetting Petal Effect in Self-Patterned Nanocolumnar Films on Poly(dimethylsiloxane)

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Abstract: Switchable mechanically induced changes in the wetting behavior of surfaces are of paramount
importance for advanced microfluidic, self-cleaning and biomedical applications. In this work we show that
the well-known polydimethylsiloxane (PDMS) elastomer develops self-patternning when it is coated with
nanostructured TiO2 films prepared by physical vapor deposition at glancing angles and subsequently subjected
to a mechanical deformation. Thus, unlike the disordered wrinkled surfaces typically created by deformation
of the bare elastomer, well-ordered and aligned micro-scaled grooves form on TiO2/PDMS after the first post-deposition bending or stretching
event. These regularly patterned surfaces can be reversibly modified by mechanical deformation,
thereby inducing a switchable and reversible wetting petal effect and the sliding of liquid
droplets. When performed in a dynamic way, this mechanical actuation produces a unique capacity of
liquid droplets (water and diiodomethane) transport and tweezing, this latter through their selective
capture and release depending on their volume and chemical characteristics. Scanning electron and atomic force microscopy studies of the strained samples showed that a dual-scale roughness,
a parallel alignment of patterned grooves and their reversible widening upon deformation,
are critical factors controlling this singular sliding behavior and the possibility to tailor their response
by the appropriate manufacturing of surface structures.

Keywords: GLAD coatings; PDMS; anisotropic wetting; droplet sliding; self-surface patterning

1. Introduction

Devices for controlling directional flows of liquids are widely utilized for self-cleaning
applications [1] or, when incorporated into microfluidic devices, to control droplet
transport processes [2]. Electrowetting [3] or mechanical actuation [4], the latter presenting
considerable limitations in terms of reversibility and long-term reproducibility, are some of the processes utilized for such a management. To promote liquid movement, wettability must be conveniently modified, acting on both the chemistry and morphology of
surfaces. Thus topographies inspired in hierarchical natural structures [5,6] and lithographic surfaces consisting of patterned lines [7,8] are commonly utilized to induce high water repellency [9] and low flow resistance [10–12]. Patterning using nanoparticles, [13]
laser [14,15], etching [16,17] or plasma treatments [18] have been also used to tailor droplet
sliding on surfaces. Similarly, a pin-release droplet effect has been reported on nanostructured thin films grown by physical vapor deposition at glancing angles (GLAD) [2]. However, those structured surfaces have a high contact angle hysteresis that precludes easy
and smooth water droplets rolling off. Conversely, oil or similar liquids with low surface free energy wet well onto such surface structures, a feature complicating the surface transport of these drops [6]. This is, for example, the case of surfaces depicting the so-called petal effect found on petal surfaces of red roses and characterized by superhydrophobicity (i.e., water contact angles higher than 150°) and a high adhesive force to water [19], both originated by a particular surface micro- and nano-structure defining a Cassie impregnating wetting state.

The effect of micro and nano-patterning on the wetting and droplet movement on flexible substrates is even less understood, despite the existence of paper-based [20] as well as vermiculite-based [21] microfluidic devices for oil–water separation, or stretchable rubber fiber microchannels [22] for conductive liquids motion and tunable pillar tip nanostructures on shape memory polymers [23]. Particularly, in elastomeric surfaces of PDMS, self-generated wrinkling processes are known to induce reversible wetting, although the phenomenon is hardly controllable due to the randomness in the generation of surface structures [24] and the dependence of water adhesion on the hysteresis contact angle [25]. Nonetheless, the fabrication of mechanically switchable wetting devices relying on anisotropically structured PDMS has been attempted with different success using complex processes such as ink transfer printing [26] and 3D printing [27], laser [28–30], bending by magnetic induction [31], wave-like nanofibers on pre-stretched substrates [32,33], incorporation of templates [34,35], nanostructures [36,37], metallic [38–40] and oxide [12,41] coatings, the selective surface functionalization through plasma treatments [42] or the deposition of superhydrophobic layers [24,43,44]. In general, these procedures lack robustness and require relatively sophisticated engineering steps to ensure straightforward functionality and full reversibility upon actuation.

Previously, we have shown that linearly ordered patterned structures self-generated on PDMS substrates covered with GLAD nanocolumnar SiO2 thin films may act as reversible optical gratings upon mechanical actuation [45]. In the present work, we propose a modified methodology based on TiO2 for the fabrication of patterned elastomeric surfaces that can be reversibly modified when subjected to strain. After a first mechanical deformation of these TiO2/PDMS samples, the self-patterned groove structures formed at their surface conferred to them an outstanding capacity to reversibly capture and release liquid droplets by mechanical actuation. This effect can be potentially useful for droplet transport and droplet tweezing capabilities. After the careful scanning electron microscopy (SEM) and atomic force microscopy (AFM) analysis of deformed specimens, we found that a key issue for this behavior is that while the pattern layout remains inalterable upon mechanical deformation, the groove width reversibly widens in some of the strained states. The relatively easy and scalable fabrication procedure of this switchable system, its stability upon long period of time and its robustness and full reversibility upon relatively unspecific mechanical actuations, sustain its use in controlled wetting, microfluidics, wearable electronics, triboelectric nanogenerators, biomechanics and optical sensors, among other advanced applications.

2. Materials and Methods

2.1. PDMS Fabrication and TiO2 Thin Film Deposition by GLAD

Polidimethylsiloxane (PDMS) foils of 2 × 2 cm² and 1.5 mm of thickness were fabricated by mixing Sylgard 184 (DOW) and two parts of silicon elastomer followed by degassing and curing at 80 °C for 30 min [45]. TiO2 thin films with 300 nm thickness were deposited on these PDMS foils in an e-beam evaporator set-up under a GLAD configuration, as previously described in [46,47] and represented in the schematic diagram of Figure S1 in the Supporting Information. A water-cooled sample holder was used to prevent radiative heating of PDMS foils during deposition of the films at a rate of 0.8–1 Å/s. Samples prepared at 60° and 85° zenithal deposition angles were called 60°- and 85°-TiO2/PDMS, respectively. For a better understanding of the reported wetting/sliding results, it is
important to consider two main directions on the surface of these TiO\textsubscript{2}/PDMS samples: one parallel of the arrival flow of material during deposition, commonly named the growing direction, and perpendicular to it, usually called the bundling direction. This latter term refers to the association of nanocolumns observed in many GLAD thin films in a direction perpendicular to the arrival flux of material [46–49] (see Supporting Information Figure S2). This lateral partial association of the TiO\textsubscript{2} nanocolumns in the TiO\textsubscript{2}/PDMS systems defines the direction of the grooves that appear upon mechanical deformation, as explained in the Results section.

2.2. Mechanical Activation and Actuation of TiO\textsubscript{2}/PDMS Samples

Regular grooves at the surface of TiO\textsubscript{2}/PDMS samples were generated when they were bent or stretched either along the growing or bundling directions, although the characteristics of the pattern (basically distance between grooves and regular distribution on the surface) were different in each case [45,49]. The TiO\textsubscript{2}/PDMS samples studied in the present work were mechanically activated along the bundling direction because the patterns were better ordered and the wetting and sliding behaviors more reliable. Once bent/stretched a first time, the samples kept memory of the generated pattern that could be reproduced so many times as the system was mechanically strained again. Characterization, wetting and drop rolling-off essays were carried out for samples in three different geometrical configurations: (1\textsuperscript{st}) flat configuration after a first bending activation, (2\textsuperscript{nd}) concave and convex bending configurations of the samples with a curvature $\kappa$ of 0.55 cm\textsuperscript{-1} and (3\textsuperscript{rd}) stretched configuration where a flat sample is stretched by a lateral elongation $\varepsilon$ ($\Delta L/L_0$, where $L$ refers to the sample length) around 0.1. The terms flat (1) concave and convex (2) and stretched (3) are kept throughout the text.

In a series of experiments, samples were dynamically stretched/bent while looking to the behavior of liquid droplets dripped and managed on their surface. In this case, the stretched sample was subjected to a larger elongation of up to 0.25.

2.3. Surface Characterization

Microstructure characterization was carried out with a Hitachi S-4800 Scanning Electronic Microscope (FESEM) at 2 kV. Surface topography was analyzed by Atomic Force Microscopy (AFM) using a Dulcinea microscope from Nanotec working in tapping mode and using high-frequency cantilevers. AFM images, taken on surface ranges of $40 \times 40$ $\mu$m\textsuperscript{2} and $5 \times 5$ $\mu$m\textsuperscript{2}, were processed with the WSxM software freely available from Nanotec [50].

Wetting characterization was performed by dripping liquid droplets on the surface of the TiO\textsubscript{2}/PDMS samples in an OCA20 contact angle system from Dataphysics. Static contact angles $\varphi$ for distilled water (WCA) and diiodomethane (DCA) were analyzed according to Young’s equation [51]. Static contact angle measurements were done with the samples in the aforementioned configurations, assuming a 3% of error that comes from the average over 3 measurements. Liquid droplets of 30 $\mu$L (water) and 5 $\mu$L (diiodomethane) were chosen for static and dynamic analysis.

Roughness, as well as structural or chemical heterogeneities, cause alterations of the ideal Young’s situation, originating a wetting contact angle hysteresis (CAH) that can be estimated through dynamic analysis, i.e., from the difference in contact angles when the volume of the deposited drop increases or decreases. These contact angles are called advancing ($\varphi_\text{a}$) and receding ($\varphi_\text{r}$), and the CAH, defined as the difference between $\varphi_\text{a}$ and $\varphi_\text{r}$, is usually related to the liquid adhesion on the surface.

Surface free energy ($\gamma_s$) was estimated by the Owens–Wendt–Kaelble method using the WCAs and DCAs [52] and taking into account the dispersive ($\gamma_d$) and polar ($\gamma_p$) components of the surface free energy according to the geometric mean method [53]:

$$\gamma_s(1 + \cos\varphi) = 2 \sqrt{\gamma_d^2 \gamma_p^2 - \gamma_d \gamma_p^2}$$  \hspace{1cm} (1)
with $\phi$ the experimental contact angle and $\gamma$ the liquid surface energy. For that estimation, we used 72.8 mJ/m$^2$ ($\gamma_{PDMS} = 50.7$ mJ/m$^2$, $\gamma_{TiO_2} = 22.1$ mJ/m$^2$) and 50.8 mJ/m$^2$ ($\gamma_{PDMS} = 2.3$ mJ/m$^2$, $\gamma_{TiO_2} = 48.5$ mJ/m$^2$) as surface energy values corresponding to water and diiodomethane, respectively.

Sliding angles ($\alpha$) were determined by tilting the sample with liquid droplets of different volume (distilled water: 2, 5, 20 and 30 µl and diiodomethane: 2 and 5 µl) deposited on its surface and looking at the angle at which they started to roll off. Measurements were taken for the three strained configurations and along two perpendicular directions: (1°) parallel to the microgrooves $\alpha^p$, and (2°) normal to the microgrooves $\alpha^n$.

3. Results

The study of the patterning in TiO$_2$/PDMS samples and their singular sliding behavior was pursued by firstly characterizing their self-arrangement behavior upon deformation, followed with a deep analysis of their wetting behavior in the light of common chemical and roughness concepts to finally accomplish the empirical determination of wetting and sliding behaviors under static and dynamic deformations in order to discuss possible deviations from current models.

3.1. Self-Patterning of TiO$_2$/PDMS Surfaces

The GLAD technique applied on flexible PDMS foils has demonstrated to give rise to TiO$_2$ thin films with oriented nanocolumnar structures. The tilting angle of the nanocolumns in the present case was $33^\circ$ and $46^\circ$ for the TiO$_2$/PDMS samples prepared by GLAD at deposition angles of $60^\circ$ ($60^\circ$-TiO$_2$/PDMS sample) and $85^\circ$ ($85^\circ$-TiO$_2$/PDMS sample), respectively (see Supporting Information Figures S1 and S2). When the “as prepared” TiO$_2$/PDMS samples were firstly deformed by bending, a new structure patterned at the microscale appeared at their surfaces, as reported in Figure 1. The distance between the parallel grooves of this structure was higher for sample $60^\circ$-TiO$_2$/PDMS (30–50 µm) than for sample $85^\circ$-TiO$_2$/PDMS (4–6 µm) and was always independent of the applied bending curvature and film thickness. When these mechanically activated TiO$_2$/PDMS samples adopted a concave configuration at a curvature of $\kappa$=0.55 cm$^{-1}$, the TiO$_2$ nanocolumns at the crack edges overlapped each other, leading to the generation of parallel protrusions. In turn, when the same mechanically activated samples adopted a convex configuration, microgrooves widened, forming angled valleys. It is worth stressing that surface transformation from concave to convex was completely reversible, and that the same patterned structures in Figure 1 were always obtained upon curving the samples as many times as desired. If instead of activating the “as prepared” TiO$_2$/PDMS samples by bending they were firstly subjected to longitudinal stretching by an elongation $\varepsilon$=0.1, parallel grooves still formed on the surface, though more separated (i.e., 17–30 µm) than when activation was carried out by bending.
Figure 1. SEM analysis of the surface state of TiO$_2$/PDMS samples subjected to different deformations after a first activation event by bending or stretching. (a) Schematic representation of the different configurations used to take the FESEM images; (b,c) FESEM micrographs of 60°-TiO$_2$/PDMS (middle) and 85°-TiO$_2$/PDMS surfaces (right). From top to bottom: (i) surface images of activated samples in a flat configuration after a first activation by bending; (ii,iii) concave and convex configurations during activation by bending and (iv) stretched samples in a flat configuration. The insets show magnified views of the groove regions. GD: average groove distance, as measured in the images.

The insets in Figure 1 showing enlarged views of the generated microgrooves reveal that in the convex and stretched states, they extend up to the PDMS substrate and that the TiO$_2$ nanocolumns become separated at the groove edges. They also suggest that the surface fraction of the exposed polymer substrate is higher for the convex than for the flat and concave configurations, being virtually zero in the latter. From SEM and AFM phase contrast images (reported as Figure S3 in the Supporting Information), average groove widths of around 225 nm and 340 nm could be determined for sample 85°-TiO$_2$/PDMS in the convex and stretched configurations, respectively. Taking into account the average groove distances (GD), determined in Figure 1, for these two situations, it could be estimated that surface area fractions of exposed PDMS amounted to 4.5% and 6.8% for the convex and stretched configurations, respectively. Similarly, the exposed PDMS area fractions for these two configurations of sample 60°-TiO$_2$/PDMS would be 2.0% and 2.3%. This means that the surface of the bent and stretched TiO$_2$/PDMS samples was heterogeneous,
likely affecting their wetting and sliding response when they were in these strained states. It is noteworthy in this regard that for the same sample and therefore similar pattern layout of grooves, these area fractions were much smaller in the flat unstrained states, a variation to which we attribute the switching behavior of the TiO2/PDMS system.

The morphological changes evidenced by the SEM analysis were accompanied by drastic changes in surface roughness. Figure 2 shows a series of AFM images taken for the PDMS substrate and 85°-TiO2/PDMS sample. Clearly, the deposition of a TiO2 nanocolumnar layer on top of PDMS substantially contributed to generate a dual-scale roughness and increase the surface heterogeneity, particularly for the three deformed configurations (see the Z scale in the AFM images). Roughness RMS values [19] in Table 1 taken for large area images (i.e., 40 × 40 μm²) confirmed a roughness enhancement in the strained samples. It is worth noting that the large zone utilized for controllable wetting analysis always encompassed one of more grooves or protrusions, thus providing a realistic view of the actual surface state.

![AFM images](image)

**Figure 2.** AFM images of PDMS and 85°-TiO2/PDMS surfaces in flat (after a single bending event, (a1,a2), stretched (d1,d2) and bent configurations, this latter in the form of either a concave (b1,b2) or a convex (c1,c2) surface, respectively.

| Roughness          | PDMS uncoated   | 60°-TiO2/PDMS | 85°-TiO2/PDMS |
|--------------------|-----------------|---------------|---------------|
|                    | RMS (nm) (40 μm × 40 μm) | RMS (nm) (5 μm × 5 μm) | RMS (nm) (40 μm × 40 μm) | RMS (nm) (5 μm × 5 μm) |
| Flat               | 17              | 1             | 44             | 15             |
| Concave            | 39              | 3             | 174            | 26             |
| Convex             | 15              | 2             | 66             | 8              |
| Stretched          | 12              | 1             | 170            | 11             |
3.2. Wetting Behavior on Strained TiO2/PDMS Surfaces

Table 2 reveals that despite the significant differences in roughness and chemical composition of the surface in contact with the liquid droplet, static water contact angles (WCAs) presented little variations for the different studied samples. Larger differences were found for diiodomethane contact angles (DCAs) that substantially decreased for the three strained configurations of samples. Furthermore, surface free energies determined according to the procedure described in the experimental section depicted a considerable increase from the flat to the strained samples with significant differences depending on sample and actual configuration. According to the data in Table 2, the dispersive component, i.e., denoting Van der Waals type interactions, was the main factor responsible for this increase in surface free energy.

**Table 2.** Wetting behavior of PDMS and TiO2/PDMS surfaces: static contact angles with water (WCA) and diiodomethane (DCA), contact angles hysteresis (CAH) in brackets and surface free energy (dispersive component in brackets) of the samples under different strained configurations.

| Surface Set-Up | WCA (CAH) [°] | DCA (CAH) [°] | Surface Free Energy (Dispersive c.) [mJ/m²] |
|----------------|---------------|---------------|---------------------------------------------|
|                | PDMS 60-TiO2 | 85-TiO2       | PDMS 60-TiO2 | 85-TiO2 | PDMS 60-TiO2 | 85-TiO2 |
| flat           | 112 (64)      | 104 (53)     | 100 (73)   | 78 (30) | 84 (17) | 80 (31) | 23 (19) | 18 (14) | 19 (15) |
| concave        | 121 (52)      | 121 (49)     | 102 (69)   | 60 (23) | 73 (33) | 63 (27) | 39 (33) | 29 (24) | 32 (27) |
| convex         | 112 (29)      | 110 (33)     | 112 (63)   | 59 (44) | 57 (44) | 54 (35) | 37 (32) | 39 (33) | 41 (35) |
| stretched      | 123 (34)      | 105 (36)     | 115 (83)   | 59 (26) | 58 (51) | 64 (24) | 40 (34) | 36 (31) | 34 (29) |

The dynamic contact angles determined for the studied samples yielded the contact angle hysteresis (CAH) values reported in Table 2. To a first approximation, CAHs could be used to predict adhesion forces and, indirectly, the sliding behavior of droplets. According to conventional descriptions [54], adhesion work (W_adh) of liquid droplets can be defined as

\[ W_{adh} = \gamma_{LV}(\cos \varphi_r - \cos \varphi_a) \]  

where \( \gamma_{LV} \) is the liquid surface free energy, and \( \varphi_r \) and \( \varphi_a \) the receding and advancing contact angles, respectively. Therefore, adhesion work increases with CAHs (calculations of this parameter are reported in Table S1 in the Supporting Information). For the set of samples studied here, the obtained values did not follow a simple tendency, except for water on samples 85-TiO2/PDMS, where this parameter was always higher than on equivalent samples of pure PDMS. However, the relatively high CAH values in Table 2 revealed that, depending on deformation state, water droplets may adhere on these surfaces while the induced variations in this parameter upon mechanical actuation forecast the possibility to control droplet sliding.

3.3. Droplet Sliding on Strained TiO2/PDMS Surfaces

Experimentally determined rolling-off angles on TiO2/PDMS were quite dependent on the actual type of sample, strained configuration and, to a lesser extent, sliding direction, either normal or parallel to the grooves. Figure 3 shows a compilation of sliding angles for water and diiodomethane droplets of, respectively, 30 µl and 5 µl, determined on 60-, 85-TiO2/PDMS and PDMS samples subjected to strain deformation. Data for sliding directions, either parallel \( \alpha' \) (black square) or perpendicular \( \alpha'' \) (blue circle) to the grooves, are also included. Water droplets presented a similar sliding angle around 35° on all the
samples in the flat configuration, while on the strained PDMS, this parameter increased to ca. 90° independently on the deformation state. A similar increase was found for 60°- and 85°-TiO2/PDMS samples in the concave configuration. However, a drastic decrease was found for the stretched (ca. 40°) and, particularly, convex configurations (15° and 25° for, respectively, 85-TiO2 and 60-TiO2/PDMS samples) in the parallel sliding direction.

![Graph showing sliding angles](image)

**Figure 3.** (a,b) Sliding angles (α|| parallel and αN perpendicular to the direction of ordered microgrooves, respectively) vs. type of strain configuration for PDMS and 60°- and 85°-TiO2/PDMS samples. The volumes of the water and diiodomethane droplets were 30 μl and 5 μl, respectively. (c) Images showing the sliding behavior of 2, 5, 10 and 30 μl water drops left on a 60°-TiO2/PDMS sample at increasing tilting angles (left), and images of a 10 μl water and 5 μl diiodomethane drops left on the 85°-TiO2/PDMS sample at increasing tilting angles (right).

For the case of non-polar liquids, the sliding behavior upon mechanical actuation depicted a high sensitivity to the direction of sliding. The sliding angles of diiodomethane were smaller than 30° for the flat configuration and slightly decreased when rolling along the grooves in comparison with the direction perpendicular to them. This sliding anisotropy was comparatively improved for the bended than the stretched configurations of the TiO2/PDMS samples. In general, stretching the samples produces a slight enhancement of the pinning state regardless of whether the liquid is polar or non-polar. However, independently of the water droplet behavior, the surface free energy of the TiO2 nanostructure causes a freed movement of diiodomethane droplets if the surface is concavely deformed. This provides a selective liquid flow for this bending. The photographs in Figure 3 together with the data in Figure S4 and video S2 in the Supporting Information also illustrate the importance of droplet weight in determining their sliding behavior; water droplets of 30 μl readily rolled-off at an angle α ≥ 30° on sample 60°-TiO2/PDMS, with a slightly favored sliding along the direction parallel to the microgrooves. It is noteworthy that despite not appreciating differences in the static contact angle along or across the groove direction, rolling-off readily occurs for water droplets of 10 μl and onwards, preferentially along such direction. This means that, for these angels and droplet volumes, gravity prevailed over adhesion interactions. Moreover, smaller diiodomethane droplets (5 μl) already slid at α < 30° on the 85°-TiO2/PDMS sample. This different sliding behavior
suggests that the TiO$_2$/PDMS system might be used for oil–water separation just by tilting or mechanical manipulation.

3.4. Dynamic Control of Droplets Sliding on Strained TiO$_2$/PDMS Surfaces

The previous experiments clearly proved a different sliding ability on mechanically deformed self-patterned TiO$_2$/PDMS surfaces in a static situation. Henceforth, we show that these stretchable materials may provide the basis to develop tunable systems for a dynamic control of water droplet sliding. Figure 4 and Video S3 from Supporting Information show two examples with stretched 60°- and 85°-TiO$_2$/PDMS samples and water and diiodomethane droplets of 30 μl and 5 μl, respectively. At that elongation (i.e., for ε~0.25, much higher than the stretching used in Figure 3), sliding of both water or diiodomethane droplets was hindered. When progressively relaxing the strain, drops started to roll off: first the diiodomethane still under a certain deformation and then water, only for the fully relaxed state. This strain dependent behavior can be qualitatively explained with the schemes in Figure 4. In the fully stretched sample, the widened cracks would produce a significant increase in the relative area of exposed PDMS in contact with the liquid and a response similar to that of this material (i.e., a sliding angle close to 90°, c.f. Figure 3). Then, when decreasing the strain and therefore the relative area of exposed PDMS, an activity identical to that in Figure 3 for unstrained samples can be obtained. Similar switchable sliding was observed when bending the samples, although experiments were less conclusive because of the practical difficulty of manually avoiding the lateral sliding of water droplets (see Supporting Information Figure S5).

Figure 4. Rolling-off control of liquid droplets on TiO$_2$/PDMS samples. (a) Photographs of 30 μl water droplet on a 85°-TiO$_2$/PDMS sample in a highly stretched state (top), which is progressively released (bottom); (b) idem for 5 μl diiodomethane droplets on 60°-TiO$_2$/PDMS sample. Green
arrows indicate the direction and magnitude of stretching. Red arrows indicate the sequence of images. The schemes at the left and right sides represent the self-patterned surface for the samples in the stretched (top) and non-stretched (bottom) states, highlighting the different distribution of PDMS (dark regions in the grooves) and TiO$_2$ (light regions) zones in the examined samples. Positive sliding is represented by green circles and pinned situations by red circles.

Overall, the experiments reported in Figures 3 and 4 clearly confirm that a controlled and switchable sliding of water and diiodomethane droplets is possible on patterned TiO$_2$/PDMS surfaces. We also learned that the observed outstanding behavior was the result of the interplay between relatively complex phenomena involving the formation of self-organized ordered grooves, linked to the bundling of TiO$_2$ nanocolumns in a preferential direction, the contribution of surface roughness at two length scales and the templated chemical heterogeneity of the surface.

4. Discussion

The experiments and results described above have been carried out with a flexible PDMS foil covered by a nanostructured TiO$_2$ thin film prepared by evaporation at glancing angles. These thin films depict a nanocolumnar structure equivalent to that obtained on rigid substrates [46] and very similar to that previously reported for SiO$_2$/PDMS systems prepared by a similar procedure [45]. Both in this previous work with SiO$_2$ and in the present paper with TiO$_2$, the nanocolumnar association character of the deposited thin film and the possibility that it offers to accommodate local stresses at the foil interface seems to be the critical factors leading to the generation of grooves upon mechanical deformation. The results herein show that the appearance of such a partially ordered groove structure is critical for the efficient and reversible control of a gradually anisotropic wetting and liquid sliding behavior on the surface of the mechanically actuated TiO$_2$/PDMS system. Key features for this behavior are the appearance of a dual scale roughness of oriented nanostructures at the surface and the small average separation between patterned grooves and protrusions, the latter meaning that one or more of these features are covered by the small liquid droplets used for the experiments. Precisely, the heterogeneous character of these surfaces, exposing PDMS at the grooves and nanocolumnar TiO$_2$ in the other zones, and the possibility to mechanically modulate the relative surface free energy contributions of these zones, are the main factors enabling an accurate mechanical control of the liquid–solid interaction in the strained and compressed states of the surface.

The wetting behavior of PDMS alone is known to be affected by mechanical deformation [55]. This capacity has been generally linked with rearrangement processes of surface groups affecting the partition between methyl groups and the more polar siloxane groups (Si-O), the latter being the majority in the flat sample. Rather than to this intrinsic behavior of PDMS, we attribute the singular behavior of the TiO$_2$/PDMS surfaces to its mechanically controlled heterogeneous character. In this regard, the higher polar component of surface free energy found for all TiO$_2$/PDMS samples with respect to PDMS can be tentatively attributed to the Ti-O (and Ti-OH) polar terminations of TiO$_2$ nanocolumns.

As a direct consequence, the rolling-off of droplets appears to be mechanically controlled on the self-patterned TiO$_2$/PDMS surfaces. According to classical sliding models, the rolling-off angle \( \alpha_{th} \) can be accounted for by the following equation:

\[
\sin \alpha_{th} = \frac{W}{mg} \gamma_{LV} (\cos \varphi_r - \cos \varphi_a)
\]

(3)

where \( m \) is the mass of the droplet, \( g \) the gravity constant, \( w \) the maximum width of the contact area of the liquid droplet and the surface and \( \gamma_{LV} \) the liquid surface free energy. The difference between the cosine of advancing and receding contact angles is related to the adhesion work \( W_{ad} \) between droplet and surface. Regarding this expression (3), the lack of clear tendencies for the CAHs (Table 2) precludes its use to predict the sliding of liquid droplets onto TiO$_2$/PDMS surfaces, a feature already reported for patterned systems [56] and attributed to inherent pinning and dewetting effects on these type of surfaces.
Another utilized parameter to predict sliding on complex surfaces is the difference between the maximum and minimum contact angles ($\phi_{\max}$ and $\phi_{\min}$), taken at the leading and trailing edges for a deformed drop just before sliding on an inclined surface [58]. Measurements of these parameters (see Table S2 from Supporting Information) did not predict completely the experimentally determined sliding angles for the different TiO$_2$/PDMS surfaces subjected to various mechanical deformations. We assume that the main factor controlling sliding and droplet movement, overcoming the hysteresis and promoted by gravitational contributions, is a wettability driving force affected by the oriented dual-scaled surface microstructure instead of the surface free energy [26].

According to our results, while on the strained PDMS the sliding parameter increased to ca. 90° independently on the deformation state [59], on the self-patterned TiO$_2$/PDMS samples, sliding depends on the strain type. Qualitatively, a simple explanation inspired by Equation (3) accounts well for the increase in rolling-off angle found for TiO$_2$/PDMS samples in the concave configuration, a result that could be associated to the increase in $\omega$ resulting from the high roughness of the developed surface in this case (Table 1) [60]. It is expected that these mechano-actuated surfaces will depict accentuated pinning effects derived from an enhanced dual-scaled roughness (Figures 1 and 2), a higher density of microgrooves and the overlapping between the bundled nanocolumns of TiO$_2$ resulting from this geometrical deformation (i.e., TiO$_2$: bundled nanocolumns at the two sides of the groves may touch and “close” the PDMS surface under the concave configuration).

However, such an explanation would not apply to the convex configuration, where the drastic decrease in sliding angle should be discussed in the context of its surface heterogeneity and the preferential orientation of the grooves generated at the surface. Moreover, since the rolling-off angle in the direction perpendicular to the grooves was slightly higher (c.f. Figure 3) than in the opposite direction, we can conclude that the preferentially oriented microstructure of the surface contributes to guide the movement of the drops through a decrease in the sliding angle. Previous works in the literature reporting that water droplets slide better along the direction of oriented microstructures have accounted for this behavior in terms of the preservation in the continuity of the three-phase (solid–water–air) contact line (TLC) [10,61]. A comparative analysis of the sliding angles for the patterned surfaces in the convex 60°-TiO$_2$/PDMS and 85°-TiO$_2$/PDMS samples supports this view, contrary to what has been observed previously in rice leaf-like wavy surfaces [5]. The inter-groove distances in these samples were, respectively, 40 µm and 6 µm, meaning that roughly 88 and 550 microgrooves are in contact with the water droplets used for the analysis and that sliding should be more favorable in the latter case. It is also noteworthy that the sliding angle along the direction perpendicular to the grooves is higher for the 85°-TiO$_2$/PDMS sample, in agreement with the loss of a directionality effect in this sample characterized by a high exposed ratio of PDMS (i.e., 2.0% vs. 4.5% for the 60°-TiO$_2$/PDMS and 85°-TiO$_2$/PDMS samples, respectively) (see Video S1 from Supporting Information). In the strained samples, the influence of differentiated TiO$_2$ and PDMS zones on sliding is further confirmed by the relatively higher sliding angles found for the stretched samples (i.e., 2.3% and 6.8% of exposed PDMS for 60- and 85-TiO$_2$/PDMS samples, respectively), where the lower surface density of microgrooves hampers a preeminent role of directionality for the control of droplet sliding.

These considerations support the possibility to achieve a precise control of the liquid movement through strain application, a concept that has been demonstrated previously for water droplet movement and retention [30]. In comparison with similar water adhesive–repulsive responses reported by Wong et al. [32] on switchable surfaces modified through the deposition of wave-like nanofibers on pre-stretched PDMS substrates, the TiO$_2$/PDMS samples depict a better-defined pattern structure that only depends on the (highly reproducible) nanostructural characteristics of the TiO$_2$: thin film prepared by GLAD and not on the stretching state of PDMS during deposition. Moreover, in agreement with the results reported by Mazaltarin et al. [62] on superhydrophobic patterned PDMS substrates, the favorable sliding found here with a marked anisotropy for the
patterned nanocolumnar TiO2/PDMS surfaces under applied strain in a convex configuration can be described as a kind of stretching deformation modified by a negative curvature.

Diodomethane, taken as example of a low surface tension liquid, presents an easier sliding behavior than water on the different studied surfaces. Moreover, the wetting by a non-polar liquid under strain is more directional, as evidenced by a notable anisotropic sliding that depended on whether the tension was applied in a direction parallel or perpendicular to the surface pattern. Thus, in accordance with the considerations above about the TLC, diiodomethane sliding angles slightly decreased when rolling along the groove direction in comparison with the sliding in a direction perpendicular to them. The different sliding behavior between water and low surface tension liquids suggests that the TiO2/PDMS system is quite appropriate as a mechanically controlled oil–water separation actuator, regulated by the interplay of surface free energy, adhesion work and gravity effects resulting from a dual scale surface roughness self-organization.

Overall, this original behavior on the TiO2/PDMS surfaces gives rise to a special droplet holding/releasing mechanism that could be taken as a mechanically controllable petal effect [19]. This behavior is different from the known transition from the lotus to the petal effects under application of a critical strain in superhydrophobic elastomeric templates [30]. The basis for the petal-like behavior depicted by our system resides in the management of the dual roughness reached by the combination of air-filled porous TiO2 nanocolumns with patterned PDMS microgrooves, giving as a result a Cassie impregnating wetting surface [62], adjustable by mechanical actuation. This mechanically tunable petal wettabillity would be useful for the transport of water droplets without the need to apply large strain deformations, as in other stretchable superhydrophobic surfaces reported previously [26,62].

5. Conclusions

In this work we have proved that the mechanical deformation of flexible samples consisting of oriented nanocolumnar TiO2 thin films deposited on PDMS produces the spontaneous formation of an ordered pattern of parallel grooves on their surface. An outstanding characteristic of these samples was that their singular patterned surface morphology, characterized by a stable groove arrangement and dimensions, can be reversibly modified by successive deformations, either bending or stretching. In addition to describing this self-organized behavior, we have demonstrated that these surface modifications can be used to reversibly control liquid wetting and droplet sliding under mechanical actuation. The fact that quite different responses were found for water and diiodomethane droplets supports the use of this type of TiO2/PDMS samples for separation, tweezing and transport of liquids droplets, acting as switchable systems where dynamic deformation (bending and stretching) could selectively delay and stop the motion of polar and non-polar liquids. This controllable wetting and sliding response of the micropatterned TiO2/PDMS surfaces is reminiscent of the typical behavior of a switchable petal effect system and sustains the fabrication of active surfaces with high prospects of use for the development of selective microfluidic devices, self-cleaning surfaces and water/oil separation purposes, among others.

Supplementary Materials: The following are available online at www.mdpi.com/2079-4991/11/10/2566/s1, Figure S1: Schematic diagram of preparation process of the self-patterned nanocolumnar TiO2 surfaces on PDMS foils; Figure S2: Cross section SEM micrographs of TiO2 nanocolumnar films deposited at 60° (a) and 85° (b) of zenithal angle. Scheme of the growing and bundling directions according to the OAD geometry; Figure S3: Phase contrast AFM images under convex bending and stretching to identify the microgroove region of exposed PDMS on the surface of the 85-TiO2/PDMS sample; Figure S4: Water contact angles (WCA) and sliding angles along the parallel (وب، اب) and normal (وب، اب) directions of the microgrooves generated onto the nanocolumnar 60-TiO2/PDMS sample in function of the droplet volume. Corresponding values onto a reference PDMS foil for a 30µl water droplet are included for comparison Figure S5: Selective water adhesion
on PDMS, 60°-TiO₂/PDMS and 85°-TiO₂/PDMS surfaces in a convex bent configuration. Pictures report the behavior on bent configuration increasing the curvature from top to bottom in each couple of panels. Drop volume was 30 μL; Table S1: Adhesion work of droplets on the different surface configurations for the TiO₂/PDMS samples and on the PDMS substrate; Table S2: Theoretical sliding angle (α) for a 30 μL water droplet from the experimental minimum (ϕ_min) and maximum (ϕ_max) contact angles values previous to rolling off (see image on the left). Data correspond to a flat configuration under stretching for the PDMS substrate and the 60°- and 85°-TiO₂/PDMS surfaces. Video S1: anisotropic wetting of 60°- and 85°-TiO₂/PDMS surfaces when subjected to convex bending in the parallel or perpendicular directions to the patterned microgrooves. Video S2: selective liquid droplet rolling-off on the 85°-TiO₂/PDMS surface under increasing tilting angles. Video S3: control of water and diiodomethane droplets sliding by stretching of 60°- and 85°-TiO₂/PDMS surfaces on a tilted configuration.

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