New memory devices based on the proton transfer process

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Abstract

Memory devices operating due to the fast proton transfer (PT) process are proposed by the means of first-principles calculations. Writing information is performed using the electrostatic potential of scanning tunneling microscopy (STM). Reading information is based on the effect of the local magnetization induced at the zigzag graphene nanoribbon (Z-GNR) edge—saturated with oxygen or the hydroxy group—and can be realized with the use of giant magnetoresistance (GMR), a magnetic tunnel junction or spin-transfer torque devices. The energetic barriers for the hop forward and backward processes can be tuned by the distance and potential of the STM tip; this thus enables us to tailor the non-volatile logic states. The proposed system enables very dense packing of the logic cells and could be used in random access and flash memory devices.

Keywords: graphene nanoribbons, memory writing, magnetism, proton transfer

(Some figures may appear in colour only in the online journal)

1. Introduction

Memory devices require densely packed systems, operating with high speed and involving stable processes for non-volatility. They also should not dissipate much heat—for energy saving reasons. Various mechanisms are used for switching the logic states in memory devices. A review on magnetic memory devices can be found in the Noble lecture by Fert [1]. Merged magnetic and superconducting switching devices have been proposed recently [2]. Organic resistive memory devices with a high on/off ratio are also available [3], and these systems can be 3D integrated [4, 5]. Other 3D integrated systems can be constructed for the optical memory devices based on photochromic molecules [6]. The basic idea of the devices proposed in this work is to use the electric field for writing the logic state and the local magnetization for reading information.

It is already known, that in the zigzag graphene nanoribbon (Z-GNR) with edges saturated with different species, the local magnetic moment arises [7, 8]. The above phenomenon was already planned to be used in spintronic devices for spin-current filtering [9, 10]. Depending on a choice of the species used for the saturation and their alternation along the edge, the local moments show up at both the edges or only at one of them. The whole edge might be magnetic or only some segments show nonvanishing local moments. The effect is very close to the edge or extends over a few bond lengths. If the local magnetization is well pronounced, even for a frequent alternation of different species along the edge—or when the segments are short, i.e. the logic cells are small—then the highly packed memory devices can be built. Using the above phenomenon in connection with a giant magnetoresistance (GMR)-based detector [11] enables very fast reading of information.

The information encrypted at the Z-GNR edge depends on its saturation. Therefore, one can saturate edges with oxygen atoms and OH groups since they adsorb via double and single bonds, respectively. Thus, writing information could be attained with the use of the proton transfer (PT) phenomenon. The hydrogen atom can move from the edge saturated with OH to the scanning tunneling microscopy (STM) tip, which is decorated with some small aromatic molecules saturated with O; and vice versa (from the STM tip to the Z-GNR). Binding/releasing the hydrogen atom to/from the molecule attached to the STM apex is similar to the vertical and horizontal transport of small molecules in the surface-engineering process [12, 13], with the difference that the molecular transport occurs directly from the metallic tip, while in our case, the molecule is always attached to the tip and only H jumps.
The PT process between the aromatic molecules (or the graphene edge) is fast—a few picoseconds or even less [14–16]. The total technological process, however, would be restricted by the speed of the alternation of the electric field—which is in the order of a gigahertz in scanning electron microscopy. On the other hand, this part could be replaced by new terahertz technology. The energy barriers for the jump forward and backward are tunable by the distance and potential of the STM tip. Hence, the high speed and non-volatile memory cells could be designed. The processes involved in the proposed system did not require an electric current flow. Therefore, the device would not be overheated and the energy dissipation would be low, as well as the scaling limits (characteristic for the resistive memories) would not be a problem [17, 18].

The above writing mechanism, namely the PT within the O-H..O and O-H...N groups, is a very popular process in organic systems. This is the basic mechanism responsible for the helical structure of DNA double strands and the secondary structure of proteins. The molecules subjecting to the PT phenomenon have been proposed as building blocks of electronic devices, such as rectifiers and transistors [19], photo and electronic devices, such as rectifiers and transistors [19], photo and optical data storage [20, 21] and optical data storage [20, 21] and optical data storage [22, 23, 24].

The binding or releasing of the hydrogen atom, in organic systems, induces a change of the single and double bond pattern, which can be mapped in the STM images. The PT process is also correlated with a change of the dipole moment—which, however, cannot be directly measured. Recent progress in material science is fast. Thus, the molecular prototypes based on the PT process might be utilized in memory devices similar to the Z-GNR edge, in the near future.

In this work, new memory devices are theoretically proposed by means of density functional calculations, performed to describe the effects involved in the reading and writing of information. The subsequent section, devoted to the results, is organized as follows. First, a presentation of the molecules, involved in the origin of our idea, and their properties is undertaken. Next, the energetic characterization of the PT in some model systems is shown. In the third part, the analysis for various cases of the Z-GNR edge saturation and the corresponding local magnetizations induced at the edges is performed. At the end of the results section, the technical details of the calculations are given.

2. Results

2.1. Molecules with PT

The PT in the molecular systems is schematically presented in figure 1(a) for two geometric configurations (A and B), which differ from each other by the location of the ‘mobile’ hydrogen atoms. These configurations are called ‘1-unit long’, because one proton-hop O-H...O pair exists symmetrically on the two sides of the axis parallel to the permanent dipole moment of the molecule. The two configurations (A and B) possess dipole moments opposite to each other. An applied electric field additionally polarizes the molecules and induces the corresponding component of the dipole moment. Figure 1(b) shows the induced dipole moment for two configurations (A and B) of the 2-unit long molecules versus the applied voltage. These molecules could be used as memory units, if it was possible to detect the dipole moment in a direct way.

The organic systems and honeycomb carbon structures, such as graphene and the carbon nanotubes (CNTs), are characterized by a network of alternated single and double bonds. This alternation also characterizes the configurations A and B (in the opposite phase)—oxygen is connected by the double bond and the OH group by the single bond. Therefore, the PT induces changes in the alternation of the bond order. Symmetric systems, it means the swapping of the configurations along the axis parallel to the dipole moment. From the top view, the configurations A and B are terminated with the horizontal double or single bond branch, respectively. These details can be seen in the STM images; such maps are presented in figure 1(c) for the 2-units long molecules. The STM imaging is, however, not a good method to detect the type of the PT configuration, because the applied electrostatic potential might induce the PT process; the measured logic state is switched, depending on the sign of the STM tip potential and the height of the energetic barrier for the PT.

The PT molecules, 2 units long, attached to the CNT (6,0) are presented in figure 2(a) for two configurations, A
and B. These cases are characterized by different I-V curves, reported in figure 2(b). The current gear of the two configurations is, however, not pronounced well enough to be used in the logic devices operating at low-bias conditions—the high voltage will change the measured logic state, in the same way as the high potential of the STM tip.

2.2. Energy barriers and writing the logic state

The total energies of the conformers, A and B, of the symmetric molecules are the same in the absence of the external electric field. In contrast, in the applied electric field, the two configurations of the 2-unit long molecules are energetically distinguished; figure 3 presents the energy values relative to the total energy at the zero field. The configuration whose dipole moment is parallel to the field vector has lower energy. The energy difference between the two configurations grows with the electric field, up to some critical value of the E-field (with the maximum of the energy separation between A and B). Additionally, the E-field polarizes the molecules, which is the source of nonlinear effects at higher values. The energies of the two configurations become closer again and swap by the successive transfer of protons at the second critical value of the E-field—this is the next zero-energy crossing point.

Detection of the logic states of the molecules—switched by the PT—is possible indirectly via the effect of the local magnetization, induced in the system with the frustrated network of the single and double bonds. Such frustration can be generated in the Z-GNR, which is saturated differently at opposite edges, or there is an alternation of the chemical termination along the same edge. The edges might end for example with oxygen or the OH group. A change of the edge saturation might be realized via the PT from the edge ended with the OH group to the STM tip apex decorated with a small aromatic molecule with one H replaced by O; and vice versa. Recently, it has been possible to move a small molecule from one site at the surface to some other adsorption site, using the STM tip. This can be done by changing the potential during the vertical and horizontal movements [12, 13]. In the same way, it is also possible to add or remove hydrogen to/from the graphene edge.

The energy barriers for such PT process between the Z-GNR edge and the STM tip molecule are schematically printed in figure 4(a). The left-hand side molecule represents a part of the Z-GNR with the edge saturated with the OH group. The right-hand side represents an aromatic molecule attached to the STM tip, and it binds the oxygen atom which is able to capture protons. In the electric field oriented towards the STM tip, the total energy of this model system is lowered when hydrogen is bound to the molecule at the STM tip (right side), and not to the Z-GNR edge (left side). The energetic barrier to be conquered for the hop forward is defined as the energy difference between the transition state configuration—which is the one with the proton in the middle of the way between the Z-GNR and the STM tip (the top configuration)—and the initial state, with the proton bound to the Z-GNR edge (left side). The barrier for the hop backwards is the energy difference between the transition state and the final state (right side). It is higher than the barrier for the forward process, which makes the system prefer the one-way reaction.
Tuning the hop-forward and hop-backwards barriers is possible via changing the distance between the Z-GNR edge and the STM tip apex—reported in figure 4(b) for the vanishing electric field—and also by varying the STM tip potential. For non-volatility, it is better to choose the larger distance between the Z-GNR edge and the STM tip apex and the higher electric field, rather than the closer distance and the smaller field. This is because the distance increases the barriers between various configurations, and the electric field increases the difference between the initial and final configuration (thus, also between the hop-forward and hop-backward barriers).

Writing in the memory device could operate as follows: we start with the STM tip possessing the OH group at the attached molecule, and we point it to the site at the Z-GNR edge which is saturated with O. Switching the potential at the tip to the positive value will move H to the Z-GNR. The molecule at the tip is now saturated with pure O. We keep the positive potential at the tip, such that none of the hydrogens from the Z-GNR can jump to the tip. We move the tip to the Z-GNR site from which we want to take H, and we switch the tip potential to the negative value—such that, the H atom jumps to the tip. Now, the STM tip state is the same as it was at the start, and it can be used again to add the H atom to the Z-GNR site which is saturated with O. We do it by moving the tip and keeping the negative potential on the way, and we change the potential to the positive value when the tip points towards the desired O atom at the Z-GNR.

Since the STM device operates on the samples in a vacuum, the well chosen distances and potentials should give enough guaranty for the correct coding. In every moment, the tip can be used for only one type of the operation: either adding or removing the H atom to/from the Z-GNR edge. Therefore, one needs to move the tip along the edge forward and backward—when it is necessary—keeping the potential set to the value, which prevents the H jump (i.e. such that the hopping barrier is higher), while changing the tip’s position. The current state of the tip is always dependent on the previous operation. The time for the PT from the tip is similar to that of moving H to the tip. This is because the aromatic molecule at the STM tip apex is similar to the graphene edge, from the saturated atom point of view.

The impractical point to consider is that STM machines are big devices. On the other hand, such a device could be simplified by removing the part which is responsible for making the images of the surface—since we only need to precisely move the tip and change the potential. The future of these devices might focus on downsizing.

2.3. Magnetic reading from the memory units

Kondo magnetism in graphene with defects is well established [25]. The nature of the defected ‘bulk’ nanoribbon states in comparison to the metallic edge-states has also been studied [26]. The effect that is the focus of this work, namely the Z-GNR edge magnetization, has been discovered by Fujita et al [7]. Recently, it has been proposed to use this phenomenon for spin filtering the current, flowing along the Z-GNR junction with differently saturated edges [9, 10].

Figure 5(a) schematically shows the mechanism of the Z-GNR edge magnetization. When both edges are saturated via single bonds, then the double bonds in the network alternate the single bonds along the zigzags parallel to the edges. In the reversed situation, when both edges are saturated via double bonds, then the double bonds in the Z-GNR network alternate the single bonds along the zigzags perpendicular to the edges. Different kinds of saturation at the opposite edges, or a change of the chemical termination along the same edge, introduces a frustration in the Z-GNR. Therefore, the electrons, which should form double bonds at the edge, are released from the chemical bonds due to the above effect of the frustration, and they are involved in the local magnetization.

The spin maps of various edge-saturation cases are presented in figures 5(b)–(g). In the homogeneously saturated Z-GNR with one edge ended with O and the opposite one with the OH group—in figure 5(b)—the ferrimagnetic state forms. It characterizes with the magnetic moments localized at the oxygens and a kind of ‘spin-wave’ at the OH-edge. The total magnetic moment per elementary cell (a cell with 19 atoms in a zigzag across the graphene nanoribbon width) is...
0.28 \mu_B and the absolute magnetic moment is 0.47 \mu_B. If the above case is modified such that the O-saturated edge is alternated with OH saturation every four units—as presented in figure 5(e) for the 8-units cell—then the magnetic moment at the O-edge segments is quenched. Although, the total and absolute magnetic moments per unit are similar to the earlier case—since now these values are 2.26 \mu_B and 3.95 \mu_B, respectively, per 8-units long cell. When the first case saturation with O and OH at opposite edges is modified by a ‘flip’ of the edge-saturation in every fourth unit—as it is in figure 5(c)—then only the O-saturated ends possess the local moments. These local moments at the opposite edges couple to each other antiferromagnetically—with a very small energetic gain, of 1 meV, with respect to the FM coupling. The absolute moment per 8-units cell is 2.27 \mu_B in both the AF and FM cases, and the total moment in the FM case is 1.76 \mu_B.

It seems that in all cases, the single and double bond alternation along the edge is never frustrated. If one edge is totally O-saturated and the opposite edge saturation is alternated with O and OH at every fourth unit, then only one edge (that with the O-ended parts between the OH-ended segments) shows the local magnetic moment and only at oxygens—as shown in figure 5(d). The total and absolute magnetizations per 8-unit long cell in this case are 0.77 \mu_B and 1.01 \mu_B, respectively. The Löwdin population analysis shows a little bit of a smaller charge at the magnetized oxygens with respect to the nonmagnetic ones, about 0.3 in the spin minority. A similar situation occurs for the alternation (at every fourth unit) of the OH–OH both edge terminated units and the

Figure 5. Edge magnetization in the Z-GNR (a) the scheme of the single and double bond network for the different edge saturations, and the spin-polarization maps, (b) ferrimagnetism in the (OH...O) case, (c) antiferromagnetic (AF)-coupling at the opposite O-edges in the (OH...O) + (O...OH) case, (d) ferromagnetic (FM) state at the (O + OH)-edge in the (O...O) + (O...OH) case, (e) ferrimagnetism at the OH-edge in the (OH...OH) + (O...OH) case, (f) AF-alternation along the (OH + O + O)-edge in the (O...OH) + (O...O) + (OH...O) case, (g) FM state at the both O-edges in the (O...O) + (OH...OH) case. \Delta n denotes the spin up and spin down density difference.
O–O both edge saturation, in figure 5(g). Oxygen on both edges possess local magnetism, with the total and absolute values of 1.8 $\mu_B$ and 2.27 $\mu_B$, respectively, per 8-units cell. When the longer segments of the oxygen saturation (8-units) reside between the OH-saturated segments (4-units) along the same edge (in the 12-unit long elementary cell) and these OH parts are not symmetrically placed on opposite edges—as indicated in figure 5(f)—the AF-coupled alternation forms along the O-segments. The total and absolute magnetizations, in this larger cell, achieve the values of 1.65 $\mu_B$ and 5.33 $\mu_B$, respectively.

Rich diversity of the magnetic orderings, existing at the edges of the Z-GNR under the inhomogeneous saturation, enables us to construct various Boolean operations. Reading the logic states could be performed using the GMR device [27]. For instance, the simplest way of reading the linear coding, would be to move the GMR device along one edge containing the segments with O and OH, while keeping the other edge completely saturated with O—this is the case shown in figure 5(d). In contrast, if the edge opposite to the side read by the GMR is totally saturated with OH groups only, there is almost no signal on the segment-alternated edge: see figure 5(e). The manipulation on both sides of the Z-GNR is more complicated and needs two GMR devices simultaneously reading the local magnetization at both edges.

Due to the fact that the system is closed in vacuum conditions, we do not have any gas reservoir of the hydrogen atoms—the number of bits 0 and bits 1 seems to be limited in the device. In order to create a source of additional H and storage space for the not needed H, we can use another Z-GNR. Thus, one Z-GNR would be an element of memory and the other should be saturated with O in half of its length and O–H in the second half. When adding or taking H from this Z-GNR reservoir, a marker of the last position for the saturation-change point is necessary. The room temperature effect for writing is very small (of about 0.025 eV at 300 K) and does not affect the energy barriers for the H transfer, which are 10 times larger.

The proposed device can operate as the magnetic random access memory, due to the high density, non-volatility and fast writing and reading processes—if one uses the GMR or magnetic tunnel junction or spin-transfer torque head for reading. Moreover, our way of writing is not limited by the large write currents in the way of Oersted-field flipping in magnetic-type writing. These currents increase the power consumption well beyond that of the static and dynamic RAM [28, 29].

For reading purposes, it is convenient to introduce separators between the bits 0 and 1. Such separators would be of the same length as a one bit and consist of the Z-GNR edge saturated with O and O–H. Thus, the reading device needs to move the head to the right bit’s position. In such a case, all bits 1 would be magnetized in the same way, independently of the length of a sequence of bits 1 and the neighboring bit; the same for bits 0. Then, reading the bits 0 would not be any problem, because the measured current is the same as in the case of zero magnetic field. To read the bit 1, one needs to use a sensitive head, which is able to detect the current difference at the magnetic field of 1 Bohr magneton at room temperature. In order to deal with a decrease of the tunnel magnetoresistance current at elevated temperatures, one could use, for instance, the Heusler compounds as reading heads [28].

2.4. Theoretical details

The molecular calculations for the energy barriers in the external electric field were performed with the quantum chemistry package TURBOMOLE [30]. The wavefunctions are represented, in the code, in the localized basis set; our choice was the correlation-consistent valence double-zeta Gaussian basis set with the polarization functions for all atoms (cc-pVDZ) [31]. The computational method was the density functional theory (DFT) [32] in the generalized gradient approximation scheme. The functional has been chosen for B3LYP [33], which is a combination of the DFT and 20% of the exact exchange—necessary for a qualitatively accurate description of the hydrogen bonds.

The calculations of the graphene-edge magnetization were performed with the QUANTUM ESPRESSO package [34]. These codes use the plane-wave basis set and pseudo-potentials to describe the core electrons. The exchange-correlation functional was chosen for the Perdew–Burke–Ernzenhoff parametrization [35]. The ultrasoft pseudopotentials have been used, with the energy cutoffs of 30 Ry and 240 Ry for the plane-waves and the density, respectively. For the edge magnetization calculations in the Z-GNR, the Monkhorst-Pack uniform k-mesh in the Brillouin zone has been set to $80 \times 1 \times 1$ for the 1-unit cell, $40 \times 1 \times 1$ for the 8-units cell, and $20 \times 1 \times 1$ for the 12-unit cell. The Fermi-surface energy broadening parameter of 0.005 Ry has been chosen for a better convergence. The vacuum separation between the periodic slabs was 20 Å.

The dipole moment and the current–voltage curves have been obtained with the Wannier functions post-processing tools [36]. Firstly, the DFT band-structure calculations were performed. Further, the maximally-localized Wannier functions have been obtained [37, 38]. In the end, the dipole moment has been calculated according to the modern theory of polarization [39, 40]. The ballistic transport calculations were performed according to the Landauer–Büttiker approach using the transport tool [41] within the Wannier90 suite of codes.

The spin-polarization maps and the STM image have been prepared using the XCrySDen package [42].

3. Conclusions

In this paper, new random access and flash memory devices are proposed by means of first-principles predictions. The basic concept relies on the PT from the OH-saturated Z-GNR edge into the STM tip decorated with a molecule, which binds O, and vice versa. Such processes are very fast, and their
energetics can be tuned by the distance between the tip and the Z-GNR edge, as well as the electrostatic potential; this thus enables picosecond and non-volatile writing of the information to the logic cells.

The PT process reverts the local dipole moment and changes the pattern of the single and double bond network in the organic and carbon honeycomb systems. This mismatch between the variously patterned networks of the chemical bonds, in turn, causes the local magnetic moment in the Z-GNR with the edges saturated differently on the opposite sides or along the same side. The induced magnetic moment can be read quickly by the GMR device.

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References

[1] Fert A 2008 Nobel lecture: origin, development, and future of spintronics Rev. Mod. Phys. 80 1517–30

[2] Baek B, Rippard W H, Benz S P, Russek S E and Dresselhaus P D 2014 Hybrid superconducting-magnetic memory device using competing order parameters Nat. Commun. 5 3888

[3] Cho B, Kim T-W, Song S, Ji Y, Jo M, Hwang H, Jang G-Y and Lee T 2010 Rewritable switching of one diode-resistor nonvolatile organic memory devices Adv. Mater. 22 1228–32

[4] Song S, Cho B, Kim T-W, Ji Y, Jo M, Wang G, Choe M, Kahng J H, Hwang H and Lee T 2010 Three-dimensional integration of organic resistive memory devices Adv. Mater. 22 5048–52

[5] Cho B, Song S, Ji Y, Kim T-W and Lee T 2011 Organic resistive memory devices: performance enhancement, integration, and advanced architectures Adv. Funct. Mater. 21 2806–29

[6] Ogawa K 2014 Two-photon absorbing molecules as potential materials for 3D optical memory Appl. Sci. 4 1–8

[7] Fujita M, Wakabayashi K, Nakada K and Kusakabe K 1996 Peculiar localized state at zigzag graphite edge J. Phys. Soc. Jpn. 65 1920–3

[8] Son Y W, Cohen M L and Louie S G 2006 Half-metallic graphene nanoribbons Nature 444 347–9

[9] Deng X Q, Zhang Z H, Tang G P, Fan Z Q, Zhu H and Yang C H 2014 Spin filtering and large magnetoresistance behaviors in carbon chain-zigzag graphene nanoribbon nanojunctions Phys. Lett. A 378 1540–7

[10] Deng X Q, Zhang Z H, Tang G P, Fan Z Q, Zhu H and Yang C H 2014 Edge contact dependent spin transport for n-type doping zigzag-graphene with asymmetric edge hydrogenation Sci. Rep. 4 4038

[11] Levy P M, Zhang S and Fert A 1990 Electrical conductivity of magnetic multilayered structures Phys. Rev. Lett. 65 1643–6

[12] Hla S W, Meyer G and Rieder K H 2001 Inducing single-molecule chemical reactions with a UHV-STM: a new dimension for nano-science and technology Chem. Phys. Chem. 2 361–6

[13] Meyer G, Bartels L and Rieder K-H 2001 Atom manipulation with the STM: nanostructuring, tip functionalization, and femtochemistry Comput. Mater. Sci. 20 443–50

[14] Rosi Ń, C N, Kühn 0, Manz J and Sundermann K 1998 The hydrogen-subway—a tunneling approach to intramolecular hydrogen transfer reactions controlled by ultrashort laser pulses J. Phys. Chem. A 102 9645–50

[15] Sliwa M, Mouton N, Ruckebusch C, Poisson L, Idrissi A, Aloise S, Potier L, Dubois J, Poizat O and Buntinx G 2010 Investigation of ultrafast photoinduced processes for salicylidenaniline in solution and gas phase: toward a general photo-dynamical scheme Photochem., Photobiol. Sci. 9 661–9

[16] Sekikawa T, Kobayashi T and Inabe T 1997 Femtosecond fluorescence study of the substitution effect on the proton transfer in thermochromic salicylideneaniline crystals J. Phys. Chem. A 101 64–9

[17] Amsink C J, Di Spigna N H, Nackashi D P and Franzon P D 2005 Scaling constraints in nano electronic random-access memories Nanotechnology 16 2251–60

[18] Zhuravlev V V, Meade R, Cavin R K and Sandhu G 2011 Scaling limits of resistive memories Nanotechnology 22 254027–47

[19] Fuentes M, Martín-Lasanta A, de Cifuentes L A, Ribagorda M, Parra A and Cuerva J M 2011 Organic-based molecular switches for molecular electronics Nanoscale 3 4003–14

[20] Benesch C, Rode M F, Ćirk M, Härtle R, Rubio-Pons O, Thoss M and Sobolewski A L 2009 Switching the Conductance of a Single Molecule by Photoinduced Hydrogen Transfer J. Phys. Chem. C 113 10315–8

[21] Hofmeister C, Härtle R, Rubio-Pons O, Coto, Pedro B, Sobolewski A L and Thoss M 2014 Switching the conductance of a molecular junction using a proton transfer reaction J. Mol. Model. 20 2163

[22] Sadek M, Wierzbowska M, Rode M F and Sobolewski A L 2014 Multipeak negative differential resistance from interplay between nonlinear Stark effect and double-channel current flow RSC Adv. 4 52933–9

[23] Raymo F M, Alvarado R J, Giordani S and Cajas M A 2003 Memory effects based on intermolecular photoinduced proton transfer J. Am. Chem. Soc. 125 2361–4

[24] Kawata S and Kawata Y 2000 Three-dimensional optical data storage using photochromic materials Chem. Rev. 100 1777–88

[25] Chen J-H, Cullen W G, Williams E D and Fuhrer M S 2011 Tunable kondo effect in graphene with defects Nat. Phys. 7 535–8

[26] Cantele G, Lee Y-S, Ninno D and Marzari N 2009 Spin channels in functionalized graphene nanoribbons Nano Lett. 9 3425–9

[27] Daughton J M 1997 Magnetic tunneling applied to memory J. Appl. Phys. 81 3758–63

[28] Bai Z, Shen L, Han G and Feng Y P 2012 Data Storage: Review of Heusler Compounds J. Phys. Chem. B 116 1230006

[29] Kesting A and Fullerton E E 2008 Device implications of spin-transfer torques J. Magn. and Magn. Mater. 320 1217–26

[30] TURBOMOLE V6.3 2011, a development of University of Karlsruhe and Forschungszentrum Karlsruhe GmbH, 19892007, TURBOMOLE GmbH, since 2007, https://turbomole.com

[31] Dunning T H 1989 Gaussian basis sets for use in correlated molecular calculations. I. the atoms boron through neon and hydrogen J. Chem. Phys. 90 1007
[32] Kohn W and Sham L J 1965 Self-Consistent equations including exchange and correlation effects Phys. Rev. 140 A1133
[33] Becke A D 1993 Density-functional thermochemistry. III. The role of exact exchange J. Chem. Phys. 98 5648
[34] Giannozzi P et al 2009 QUANTUM ESPRESSO: a modular and open-source software project for quantum simulations of materials J. Phys.: Condens. Matter. 21 395502
[35] Perdew J P, Burke K and Ernzerhof M 1996 Generalized gradient approximation made simple Phys. Rev. Lett. 77 3865
[36] Mostofi A A, Yates J R, Lee Y-S, Souza I, Vanderbilt D and Marzari N 2008 A tool for obtaining maximally-localised Wannier functions Comput. Phys. Commun. 178 685–15 www.wannier.org
[37] Marzari N and Vanderbilt D 1997 Maximally localized generalized Wannier functions for composite energy bands Phys. Rev. B 56 12847–19
[38] Marzari N, Mostofi A A, Yates J R, Souza I and Vanderbilt D 2012 Maximally localized Wannier functions: Theory and applications Rev. Mod. Phys. 84 1419–57
[39] Resta R 1994 Macroscopic polarization in crystalline dielectrics: The geometric phase approach Rev. Mod. Phys. 66 899
[40] Gibertini M, Pizzi G and Marzari M 2014 Engineering polar discontinuities in honeycomb lattices Nat. Commun. 5 5157–13
[41] Shelley M, Poilvert N, Mostofi A A and Marzari N 2011 Automated quantum conductance calculations using maximally-localised Wannier functions Comput. Phys. Commun. 182 2174
[42] Kokalj A 2003 Computer graphics and graphical user interfaces as tools in simulations of matter at the atomic scale Comput. Mater. Sci. 28 155