Finiteness of polygonal relative equilibria for
generalised quasi-homogeneous $n$-body problems
and $n$-body problems in spaces of constant
curvature

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Abstract

We prove for generalisations of quasi-homogeneous $n$-body problems with center of mass zero and $n$-body problems in spaces of negative constant Gaussian curvature that if the masses and rotation are fixed, there exists, for every order of the masses, at most one equivalence class of relative equilibria for which the point masses lie on a circle, as well as that there exists, for every order of the masses, at most one equivalence class of relative equilibria for which all but one of the point masses lie on a circle and rotate around the remaining point mass. The method of proof is a generalised version of a proof by J.M. Cors, G.R. Hall and G.E. Roberts on the uniqueness of co-circular central configurations for power-law potentials.

1 Introduction

By $n$-body problems we mean problems where we are tasked with deducing the dynamics of $n$ point masses described by a system of differential equations. The study of such problems has applications to various fields, including atomic physics, celestial mechanics, chemistry, crystallography, differential equations, dynamical systems, geometric mechanics, Lie groups and algebras, non-Euclidean and differential geometry, stability theory, the theory of polytopes and topology (see for example [1], [2], [15], [17], [19], [28], [53], [59], [60], [61], [66] and the references therein). The $n$-body problems that form the backbone of this paper are a generalisation of a class of quasi-homogeneous $n$-body problems, which we will call generalised $n$-body
problems for short and the $n$-body problem in spaces of constant Gaussian curvature, or curved $n$-body problem for short. By the generalised $n$-body problem we mean the problem of finding the orbits of point masses $q_1, \ldots, q_n \in \mathbb{R}^2$ and respective masses $m_1 > 0, \ldots, m_n > 0$ determined by the system of differential equations

\[ \ddot{q}_i = \sum_{j=1, j \neq i}^{n} m_j (q_j - q_i) f \left( \|q_j - q_i\| \right), \quad (1.1) \]

where $\|\cdot\|$ is the Euclidean norm, $f$ is a positive valued scalar function and $xf(x)$ is a decreasing, differentiable function. Our definition of generalised $n$-body problems thus includes a large subset of quasi-homogeneous $n$-body problems, which are problems with $f(x) = Ax^{-a} + Bx^{-b}$, where $A, B \in \mathbb{R}$ and $0 \leq a < b$, which include problems studied in fields such as celestial mechanics, crystallography, chemistry and electromagnetics (see for example \[8\]–\[13\], \[20\], \[22\], \[24\], \[28\], \[33\] and \[49\]–\[52\]).

By the $n$-body problem in spaces of constant Gaussian curvature, we mean the problem of finding the dynamics of point masses $p_1, \ldots, p_n \in \mathbb{M}_\sigma^2 = \{(x_1, x_2, x_3) \in \mathbb{R}^3 | x_1^2 + x_2^2 + \sigma x_3^2 = \sigma\}$, where $\sigma = \pm 1$ and respective masses $\hat{m}_1 > 0, \ldots, \hat{m}_n > 0$, determined by the system of differential equations

\[ \ddot{p}_i = \sum_{j=1, j \neq i}^{n} \frac{\hat{m}_j (p_j - \sigma(p_i \odot p_j)p_i)}{\left(\sigma - \sigma(p_i \odot p_j)^2\right)^{\frac{3}{2}}} - \sigma(\dot{p}_i \odot \dot{p}_i)p_i, \quad i \in \{1, \ldots, n\}. \quad (1.2) \]

where for $x, y \in \mathbb{M}_\sigma^2$ the product $\cdot \odot \cdot$ is defined as

\[ x \odot y = x_1 y_1 + x_2 y_2 + \sigma x_3 y_3. \]

The curved $n$-body problem generalises the classical, or Newtonian $n$-body problem ($f(x) = x^{-\frac{3}{2}}$ in (1.1)) to spaces of constant Gaussian curvature (i.e. spheres and hyperboloids) and goes for the two body case back to the 1830s, (see \[6\] and \[43\]), followed by \[57\], \[58\], \[36\], \[37\], \[38\], \[40\], \[41\], \[42\], \[39\], but it was not until a revolution took place with the papers \[25\], \[26\], \[27\] by Diacu, Pérez-Chavela and Santoprete in which the successful study of $n$-body problems in spaces of constant Gaussian curvature for the case that $n \geq 2$ was established. After this breakthrough, further results for the $n \geq 2$ case were then obtained in \[7\], \[14\]–\[18\], \[21\], \[23\], \[29\], \[30\] and \[62\]–\[65\]. See \[14\]–\[17\] and \[21\] for a detailed historical overview of the development of.
the curved \( n \)-body problem. In this paper we will only consider the negative constant curvature case, i.e. the case \( \sigma = -1 \).

For these two types of \( n \)-body problems we will prove results regarding the finiteness of relative equilibrium solutions of (1.1) and (1.2), which are solutions of (1.1), or (1.2), for which the configuration of the point masses stays fixed in shape and size over time. Specifically:

We will call \( q_1, \ldots, q_n \in \mathbb{R}^2 \) a relative equilibrium of (1.1) if \( q_i(t) = T(A)(Q_i - Q_M) + Q_M, \; i \in \{1, \ldots, n\} \), where \( Q_i \in \mathbb{R}^2, \; A \in \mathbb{R}_{>0} \) are constant,

\[
T(t) = \begin{pmatrix}
\cos t & -\sin t \\
\sin t & \cos t
\end{pmatrix}
\]

is a \( 2 \times 2 \) rotation matrix and

\[
Q_M = \frac{1}{M} \sum_{k=1}^{n} m_k Q_k
\]

is the center of mass with \( M = \sum_{k=1}^{n} m_k \). If the \( q_i \) lie on a circle with the origin at its center, we will call \( q_1, \ldots, q_n \) a polygonal relative equilibrium solution of (1.1). If all but one of the masses form a polygon with the origin at its center, with the remaining mass at the origin, then we will call such a relative equilibrium a polygonal relative equilibrium with center zero of (1.1) for short.

Following the example of [25], [26], [27] by Diacu, Pérez-Chavela and Santoprete, we will call \( p_1, \ldots, p_n \in \mathbb{M}^2_\sigma \) a polygonal relative equilibrium of (1.2) if

\[
p_i(t) = \begin{pmatrix}
T(Bt)P_i \\
z
\end{pmatrix},
\]

where \( P_i \in \mathbb{R}^2, \; z \in \mathbb{R}, \; B \in \mathbb{R}_{>0} \) are constant and \( i \in \{1, \ldots, n\} \). For a proof of the existence of such solutions we refer the reader to Theorem 1 of [15].

Finally, following [28], we will say that two relative equilibria of (1.1) are equivalent, or are in the same equivalence class, if they are equivalent under rotation. For the constant curvature case, we will say that two polygonal relative equilibria are equivalent if they are equivalent under a rotation induced by a rotation matrix of the type \( \begin{pmatrix}
T(c) & 0 \\
0^T & 1
\end{pmatrix} \), where \( c \in \mathbb{R} \) is a constant, \( 0 \in \mathbb{R}^2 \) is the zero vector and \( 0^T \) its transpose. It should be noted that these definitions differ from the usual definition (see for example [59]).
where two relative equilibria are also considered to be equivalent if they are equivalent under scalar multiplication.

The relevance of the relative equilibria studied in this paper is twofold: In [9], Cors, Hall and Roberts proved for the case that if $Q_M = 0$, $f(x) = x^{-\alpha-2}$, $\alpha > 0$ and if $A$ and the masses $m_1,..,m_n$ are fixed, then for every order of the masses there exists at most one equivalence class of polygonal relative equilibrium solutions (called co-circular central configurations for power-law potentials in [9]) of (1.1). This may be a significant step in the direction of proving Problem 12 of [2], an important list of open problems in the field of celestial mechanics, composed by Albouy, Cabral and Santos: Are there, except for the regular $n$-gon with equal masses, any polygonal relative equilibria of (1.1) for the case that $f(x) = x^{-a}$, $a \geq 1$? A logical step to make is to investigate to which extent Cors’, Hall’s and Roberts’ result can be applied to $n$-body problems in spaces of constant curvature, or any generalised $n$-body problems. Additionally, generalising this result may shed further light on solving Problem 12 of [2] and the sixth Smale problem, which conjectures that for any fixed set of masses, the corresponding set of equivalence classes of relative equilibria of the classical $n$-body problem is finite (see [59]). Secondly and entwined with the theoretical aspect, relative equilibria can tell us a great deal about the geometry of the universe and orbits in our solar system: It was proven in [25] and [26] that while for the zero curvature case polygonal relative equilibria shaped as equilateral triangles with unequal masses exist, in nonzero constant curvature spaces the masses have to be equal, proving that the region between the Sun, Jupiter and the Trojan asteroids has to be flat. This means that getting any information about polygonal relative equilibria that exist in spaces of positive constant curvature, zero curvature, or negative constant curvature can further our understanding about the geometry of the universe. Additionally, the ring problem, or a regular polygonal relative equilibrium with one mass at its center and all masses on the circle equal (see for example [31]) is a model that was originally formulated by Maxwell to describe the dynamics of particles orbiting Saturn (see [44]) and has since then been applied to describing other planetary rings, asteroid belts, planets orbiting stars, stellar formations, stars with an accretion ring, planetary nebula and motion of satellites (see [3], [4], [5], [31], [32], [34], [35], [45], [46], [47], [53]–[56]). In this context, considering the more general solutions of polygonal relative equilibria, proving the number of possible equilibria to be finite may be a very fruitful endeavour. We will prove the following theorems:

**Theorem 1.1.** Let $A$, $m_1,..,m_n$ be fixed. For every order of the masses,
there exists at most one equivalence class of polygonal relative equilibria of (1.1) with \( Q_M = 0 \).

**Theorem 1.2.** Let \( B, \hat{m}_1, \ldots, \hat{m}_n \) be fixed and let \( \sigma = -1 \). For every order of the masses, there exists at most one equivalence class of polygonal relative equilibria of (1.2).

**Theorem 1.3.** Let \( A, m_1, \ldots, m_n \) be fixed. Let \( n = N + 1 \). Then for every order of the masses, there exists at most one equivalence class of relative equilibria of (1.1) with \( Q_M = 0 \) where for \( i \in \{1, \ldots, N\} \) the \( q_i \) lie on a circle with the origin at its center and \( q_{N+1} = 0 \).

We will first prove Theorem 1.1 in section 2, after which we will prove Theorem 1.2 in section 3 and finally Theorem 1.3 in section 4.

## 2 Proof of Theorem 1.1

We will prove Theorem 1.1 by to a large extent following the proof of Theorem 3.2 in [9], which is reminiscent of a topological approach by Moulton (see [48], [9]), but instead of making the proof work for \( f(x) = x - \alpha - 2 \), where \( \alpha > 0 \), as was done in [9], we successfully realise the result for any positive function \( f \) for which \( x f(x) \) is a decreasing function:

If \( q_i, i \in \{1, \ldots, n\} \) is a relative equilibrium of (1.1) and \( Q_M = 0 \), we may write \( q_i(t) = T(At)Q_i \), where

\[
Q_i = r \begin{pmatrix} \cos \alpha_i \\ \sin \alpha_i \end{pmatrix}, \quad \text{for } i \in \{1, \ldots, n\}, \quad r > 0
\]

and \( 0 \leq \alpha_1 < \alpha_2 < \ldots < \alpha_n < 2\pi \), meaning that if we insert \( q_i(t) = rT(At)Q_i \) and \( q_j(t) = rT(At)Q_j \) into (1.1) and using that in that case

\[
\|Q_i - Q_j\| = r \sqrt{2 - 2 \cos (\alpha_i - \alpha_j)} = 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right)
\]

and multiplying both sides of the resulting equation with \( -T(-\alpha_i) \) from the left, we can rewrite (1.1) as

\[
r \begin{pmatrix} A^2 \\ 0 \end{pmatrix} = r \sum_{j=1, j \neq i}^{n} m_j \begin{pmatrix} 1 - \cos (\alpha_i - \alpha_j) \\ \sin (\alpha_i - \alpha_j) \end{pmatrix} f \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right),
\]

(2.1)
which, if we write \( g(x) = xf(x) \), can be rewritten as

\[
\begin{align*}
  m_i A^2 r &= \sum_{j=1, j \neq i}^n m_i m_j \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \\
  0 &= \sum_{j=1, j \neq i}^n m_i m_j r \delta_{ij} \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right),
\end{align*}
\]

(2.2)

where

\[
\delta_{ij} = \begin{cases} 
1 & \text{if } i > j \\
-1 & \text{if } i < j.
\end{cases}
\]

If \( G \) is any scalar function for which \( G'(x) = g(x) \), then defining

\[
V(r, \alpha_1, ..., \alpha_n) = \sum_{l=1}^n \sum_{k=1, k \neq l}^n m_l m_k G \left( 2r \sin \left( \frac{1}{2} |\alpha_l - \alpha_k| \right) \right) - \sum_{l=1}^n m_l A^2 r^2
\]

(2.3)

gives by (2.2) that

\[
\frac{\partial V}{\partial r} = 0 \quad \text{and} \quad \frac{\partial V}{\partial \alpha_i} = 0 \text{ for all } i \in \{1, ..., n\}.
\]

This means that, for whatever values of \( r, \alpha_1, ..., \alpha_n \) the vectors \( q_i(t) = T(At)Q_i \) give a relative equilibrium solution of (1.1), \( (r, \alpha_1, ..., \alpha_n) \) is a stationary point of \( V \). We will show that for such a stationary point \( V \) has to have a local maximum, i.e.

\[
\rho^2 \frac{\partial^2 V}{\partial r^2} + 2\rho \sum_{l=1}^n \gamma_l \frac{\partial^2 V}{\partial r \partial \alpha_l} + \sum_{l=1}^n \sum_{k=1}^n \gamma_l \gamma_k \frac{\partial^2 V}{\partial \alpha_l \partial \alpha_k} \leq 0
\]

(2.4)

for all vectors \( (\rho, \gamma_1, ..., \gamma_n) \in \mathbb{R}^{n+1} \).

Note that by (2.2)

\[
\rho^2 \frac{\partial^2 V}{\partial r^2} = 4\rho^2 \sum_{i=1}^n \sum_{j=1, j \neq i}^n m_i m_j \sin^2 \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) - 2\rho^2 \sum_{i=1}^n m_i A^2.
\]

(2.5)
which can be rewritten as

\[ 2\rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} = 4\rho \sum_{i=1}^{n} \gamma_i \frac{\partial}{\partial r} \sum_{j=1, j \neq i}^{n} m_i m_j r \delta_{ij} \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \]

\[ = 4\rho \sum_{i=1}^{n} \gamma_i \left( \sum_{j=1, j \neq i}^{n} m_i m_j r \delta_{ij} \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \right) + 2 \sum_{j=1, j \neq i}^{n} m_i m_j r \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) , \]

which, by the second identity of (2.2), gives

\[ 2\rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} = 4\rho \sum_{i=1}^{n} \gamma_i \left( 0 + 2\gamma_i \sum_{j=1, j \neq i}^{n} m_i m_j r \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \right) \]

\[ = 4\rho \sum_{i=1}^{n} \gamma_i \left( \sum_{j=1, j \neq i}^{n} (\gamma_i - \gamma_j) m_i m_j r \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \right) , \]

so combined with (2.3), this gives

\[ \rho^2 \frac{\partial^2 V}{\partial r^2} + 2\rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} = -2\rho^2 \sum_{i=1}^{n} m_i A_i^2 + \sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} m_i m_j \left( 4\rho^2 \sin^2 \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) + 4(\gamma_i - \gamma_j) r \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \]

which can be rewritten as

\[ \rho^2 \frac{\partial^2 V}{\partial r^2} + 2\rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} = -2\rho^2 \sum_{i=1}^{n} m_i A_i^2 \]

\[ + \sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} m_i m_j \left( 2\rho \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) + r(\gamma_i - \gamma_j) \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \right)^2 \]

\[ \cdot g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \]

\[ - \sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} m_i m_j (\gamma_i - \gamma_j)^2 r^2 \cos^2 \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \]

\[ = 2\rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} \]

(2.6)
Thirdly, for \( i \neq j \), by (2.2),

\[
\frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = 2 \frac{\partial}{\partial \alpha_j} \left( m_i m_j r \delta_{ij} \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \right) - 2m_i m_j r^2 \cos^2 \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right)
\]

and thus

\[
\frac{\partial^2 V}{\partial \alpha_i^2} = -2 \sum_{j=1, j \neq i}^n m_i m_j r \delta_{ij} \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) + 2 \sum_{j=1, j \neq i}^n m_i m_j r^2 \cos^2 \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right)
\]

so by (2.7),

\[
\sum_{i=1}^n \sum_{j=1}^n \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = \sum_{i=1}^n \gamma_i^2 \frac{\partial^2 V}{\partial \alpha_i^2} + \sum_{i=1}^n \sum_{j=1, j \neq i}^n \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j}
\]

\[
= - \sum_{i=1}^n \sum_{j=1, j \neq i}^n \gamma_i^2 \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} + \sum_{i=1}^n \sum_{j=1, j \neq i}^n \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = \sum_{i=1}^n \sum_{j=1, j \neq i}^n \left( -\gamma_i^2 + \gamma_i \gamma_j \right) \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j}
\]

giving

\[
\sum_{i=1}^n \sum_{j=1, j \neq i}^n \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = \sum_{i=1}^n \sum_{j=1, j \neq i}^n \left( -\gamma_i^2 + \gamma_i \gamma_j \right) \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j}
\]

\[
= \frac{1}{2} \left( \sum_{i=1}^n \sum_{j=1, j \neq i}^n \left( -\gamma_i^2 + \gamma_i \gamma_j \right) \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} + \sum_{j=1}^n \sum_{i=1, i \neq j}^n \left( -\gamma_j^2 + \gamma_j \gamma_i \right) \frac{\partial^2 V}{\partial \alpha_j \partial \alpha_i} \right)
\]

\[
= -\frac{1}{2} \sum_{i=1}^n \sum_{j=1, j \neq i}^n \left( \gamma_i^2 - 2\gamma_i \gamma_j + \gamma_j^2 \right) \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = -\frac{1}{2} \sum_{i=1}^n \sum_{j=1, j \neq i}^n \left( \gamma_i^2 - \gamma_j^2 \right) \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j}
\]
and therefore
\[ \sum_{i=1}^{n} \sum_{j=1}^{n} \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = -\sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} (\gamma_i - \gamma_j)^2 m_i m_j r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \]
\[ + \sum_{j=1, j \neq i}^{n} m_i m_j (\gamma_i - \gamma_j)^2 r^2 \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \g'(2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right). \]
\[ (2.8) \]

Combining (2.4), (2.6) and (2.8), we now get that
\[ \rho^2 \frac{\partial^2 V}{\partial r^2} + 2\rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} + \sum_{i=1}^{n} \sum_{j=1}^{n} \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} = -\rho^2 \sum_{i=1}^{n} m_i A^2 \]
\[ + \sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} m_i m_j \left( 2\rho \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) + r (\gamma_i - \gamma_j) \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \right)^2 \]
\[ \cdot \g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) + 0 \]
\[ - \sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} (\gamma_i - \gamma_j)^2 m_i m_j r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \].

As by construction \(\g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) < 0\) and \(g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) > 0\), this means that
\[ \rho^2 \frac{\partial^2 V}{\partial r^2} + \rho \sum_{i=1}^{n} \gamma_i \frac{\partial^2 V}{\partial r \partial \alpha_i} + \sum_{i=1}^{n} \sum_{j=1}^{n} \gamma_i \gamma_j \frac{\partial^2 V}{\partial \alpha_i \partial \alpha_j} \leq 0 \]
\[ (2.9) \]

with equality if and only if \(\gamma_i = \gamma_j\) and \(\rho = 0\), which can be prevented by fixing one of the \(Q_i, i \in \{1, ..., n\}\). This proves that any stationary point of \(V\) gives a maximum value of \(V\), proving Theorem 1.1.

3 Proof of Theorem 1.2

Let \(p_1, ..., p_n\) be a polygonal relative equilibrium of (1.2) and let for \(i \in \{1, ..., n\}\)
\[ P_i = \rho \left( \begin{array}{c} \cos \gamma_i \\ \sin \gamma_i \end{array} \right), \]
\[ (3.1) \]
where \( \gamma_1, ..., \gamma_n \in [0, 2\pi) \) are ordered from smallest to largest and \( \rho > 0 \). We will prove that the \( P_i \) will give rise to a system of equations in the same way the \( Q_i \) in the proof of Theorem 1.1 give rise to (2.1): Inserting (3.1) into (1.2) and multiplying both sides of the resulting equation for the first two entries of \( \dot{\tilde{p}}_i \) with \( T(-Bt) \) gives

\[
-B^2 P_i = \sum_{j=1, j \neq i}^{n} \frac{\hat{m}_j (P_j - \sigma((P_i, P_j) + \sigma z^2) P_i)}{(\sigma - \sigma((P_i, P_j) + \sigma z^2))^2} - \sigma(\dot{p}_i \odot \tilde{\dot{p}}_i) P_i, \quad i \in \{1, ..., n\},
\]

which can be rewritten as

\[
-B^2 P_i = \sum_{j=1, j \neq i}^{n} \frac{\hat{m}_j (P_j - P_i)}{(\sigma - \sigma((P_i, P_j) + \sigma z^2))^2} + \left( \sum_{j=1, j \neq i}^{n} \frac{\hat{m}_j (1 - \sigma((P_i, P_j) + \sigma z^2))}{(\sigma - \sigma((P_i, P_j) + \sigma z^2))^2} - \sigma(\dot{p}_i \odot \tilde{\dot{p}}_i) \right) P_i \quad \text{(3.2)}
\]

and the identity for the third entry of \( \ddot{p}_i \) then is

\[
0 = \sum_{j=1, j \neq i}^{n} \frac{\hat{m}_j (1 - \sigma((P_i, P_j) + \sigma z^2)) z}{(\sigma - \sigma((P_i, P_j) + \sigma z^2))^2} - \sigma(\dot{p}_i \odot \tilde{\dot{p}}_i) z. \quad \text{(3.3)}
\]

Note that \(-1 = p_i \odot p_i = \|P_i\|^2 - z^2 = \rho^2 - z^2 \) for \( \sigma = -1 \), so \( z \neq 0 \). Therefore, using (3.3), the second sum of (3.2) may be replaced with zero, giving

\[
-B^2 P_i = \sum_{j=1, j \neq i}^{n} \frac{\hat{m}_j (P_j - P_i)}{(\sigma - \sigma((P_i, P_j) + \sigma z^2))^2}, \quad \text{(3.4)}
\]

which by (3.1) can be rewritten, using that \( \sigma z^2 = \sigma - \rho^2 \) and multiplying both sides of (3.4) with \(-T(-\gamma_i)\), as

\[
B^2 \begin{pmatrix} \rho \\ 0 \end{pmatrix} = \sum_{j=1, j \neq i}^{n} \hat{m}_j \rho \left( \frac{1 - \cos(\gamma_j - \gamma_j)}{\sin(\gamma_j - \gamma_i)} \right) \\
\cdot \rho^{-3} (1 - \cos(\gamma_j - \gamma_i))^{-\frac{3}{2}} (2 - \sigma \rho^2 (1 - \cos(\gamma_j - \gamma_i)))^{-\frac{3}{2}}. \quad \text{(3.5)}
\]
As $\rho(1 - \cos(\gamma_j - \gamma_i))^{\frac{1}{2}} = \sqrt{2} \cdot \rho \sin\left(\frac{1}{2}|\gamma_j - \gamma_i|\right)$, we may rewrite (3.5) as

$$B^2(\rho) = \sum_{j=1, j\neq i}^n \hat{m}_j \rho \left(1 - \cos(\gamma_j - \gamma_i)\right) \sin(\gamma_j - \gamma_i) \left(2\rho \sin\left(\frac{1}{2}|\gamma_j - \gamma_i|\right)\right)^{-\frac{3}{2}} \cdot 8 \left(2\rho \sin\left(\frac{1}{2}|\gamma_j - \gamma_i|\right)\right)^{-3} \left(4 - \sigma \left(2\rho \sin\left(\frac{1}{2}|\gamma_j - \gamma_i|\right)\right)\right)^2 - \frac{3}{2}.$$ (3.6)

If we define $h(x) = 8x^{-3}(4 - \sigma x^2)^{-\frac{3}{2}}$, then we can rewrite (3.6) as

$$B^2(\rho) = \sum_{j=1, j\neq i}^n \hat{m}_j \rho \left(1 - \cos(\gamma_j - \gamma_i)\right) \sin(\gamma_j - \gamma_i) h\left(2\rho \sin\left(\frac{1}{2}|\gamma_j - \gamma_i|\right)\right).$$ (3.7)

Now (3.7) is an identity exactly of the same type as (2.1), as $(xh(x))' < 0$. So going through the proof of Theorem 1.1 using (3.7) instead of (2.1) now proves our theorem. It should be noted that (3.5) was already proven in [15] in a more general setting, but as the calculation is not that long for this specific case, the argument has been repeated to make the paper self contained.

4 Proof of Theorem 1.3

Let $n = N + 1$. If $q_1,...,q_n$ is a relative equilibrium of (1.1) with $Q_M = 0$, the point masses $q_1,...,q_N$ lie on a circle and $Q_{N+1} = 0$, then we may write for $i \in \{1,...,N\}$

$$Q_i = r \begin{pmatrix} \cos\alpha_i \\ \sin\alpha_i \end{pmatrix},$$

where $r > 0$ and $0 \leq \alpha_1 < ... < \alpha_N < 2\pi$. Following the proof of Theorem 1.1 again for $i \in \{1,...,N\}$, inserting these expressions for the $Q_i$ into (1.1) gives instead of (2.1) the slightly different identity

$$-A^2Q_i = \sum_{j=1, j\neq i}^N m_j (Q_j - Q_i) f(||Q_j - Q_i||) + m_{N+1}(0 - Q_i) f(||0 - Q_i||),$$

giving

$$(m_n f(||Q_i||) - A^2) Q_i = \sum_{j=1, j\neq i}^N m_j (Q_j - Q_i) f(||Q_j - Q_i||),$$
which by the same argument that gave (2.2) now gives

$$
\begin{align*}
&\begin{cases}
  m_i \left( A^2 r - m_n g(r) \right) = \sum_{j=1, j \neq i}^N m_i m_j \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \\
  0 = \sum_{j=1, j \neq i}^N m_i m_j r \delta_{ij} \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right)
\end{cases}
\end{align*}
$$

(4.1)

with again \( g(x) = xf(x) \) and \( \delta_{ij} \) as in the proof of Theorem 1.1. So as long as \( (A^2 - m_n f(r)) \) has to be positive) we can continue to go through the proof of Theorem 1.1, replacing the function \( V \) with again \( g \) for all \( r_\alpha, ... , \alpha_N \) = \( \sum_{i=1}^N \sum_{k=1, k \neq i}^N m_i m_k G \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_k| \right) \right) \) - \( \sum_{i=1}^n m_i \left( A^2 r^2 - 2m_n G(r) \right) \).

Repeating the proof of Theorem 1.1 using \( W \) instead of \( V \) leads to

$$
\rho^2 \frac{\partial^2 W}{\partial r^2} + 2\rho \sum_{i=1}^N \gamma_i \frac{\partial^2 W}{\partial r \partial \alpha_i} + \sum_{i=1}^N \sum_{j=1}^N \gamma_i \gamma_j \frac{\partial^2 W}{\partial \alpha_i \partial \alpha_j} = -2\rho^2 \sum_{i=1}^N m_i (A^2 - m_n g'(r)) \left. \right|_{r_\alpha, ... , \alpha_N}
$$

+ \( \sum_{i=1}^N \sum_{j=1, j \neq i}^N m_i m_j \left( 2\rho \sin \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) + r (\gamma_i - \gamma_j) \cos \left( \frac{1}{2} (\alpha_i - \alpha_j) \right) \right)^2 \)

- \( \sum_{i=1}^N \sum_{j=1, j \neq i}^N (\gamma_i - \gamma_j)^2 m_i m_j r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) \)

for all \( \rho, \gamma_1,...,\gamma_N \in \mathbb{R} \). As \( g' \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) < 0, g'(r) < 0 \) and \( g \left( 2r \sin \left( \frac{1}{2} |\alpha_i - \alpha_j| \right) \right) > 0 \), this means that

$$
\rho^2 \frac{\partial^2 W}{\partial r^2} + 2\rho \sum_{i=1}^N \gamma_i \frac{\partial^2 W}{\partial r \partial \alpha_i} + \sum_{i=1}^N \sum_{j=1}^N \gamma_i \gamma_j \frac{\partial^2 W}{\partial \alpha_i \partial \alpha_j} \leq 0
$$

with equality if and only if \( \rho = 0 \) and \( \gamma_i = \gamma_j \) for all \( i, j \in \{1,...,N\} \), which can be prevented by fixing one of the \( Q_i, i \in \{1,...,N\} \). We thus find that by the same argument as used in the proof of Theorem 1.1 that for any order of masses \( m_1,...,m_n \), there exists at most one relative equilibrium of (1.1) with center of mass zero, where all point masses but one lie on a circle around the remaining point mass, which coincides with the center of mass.
References

[1] R. Abraham, J. Marsden, (1978) Foundations of Mechanics, Addison-Wesley Publishing Co. Reading, Mass.

[2] A. Albouy, H.E. Cabral, A.A. Santos, Some problems on the classical n-body problem, *Celest. Mech. Dyn. Astr.* 133, (2012) 369–375.

[3] M. Arribas, A. Elipe, Bifurcations and equilibria in the extended n-body ring problem, *Mech. Res. Comm.* 31:18, (2004).

[4] M. Arribas, A. Elipe, T. Kalvouridis, M. Palacios, Homographic solutions in the planar n + 1-body problem with quasi-homogeneous potentials, *Celest. Mech. and Dyn. Astr.*, 99(1), (2007) 1-12.

[5] M. Arribas, A. Elipe, M. Palacios, Linear stability of ring systems with generalized central forces, *Astron. Astrophys.*, 489, (2008) 819-824.

[6] W. Bolyai and J. Bolyai, *Geometrische Untersuchungen*, Hrsg. P. Stäckel, Teubner, Leipzig-Berlin, 1913.

[7] J.F. Cariñena, M.F. Rañada and M. Santander, Central potentials on spaces of constant curvature: The Kepler problem on the two-dimensional sphere $S^2$ and the hyperbolic plane $H^2$, *J. Math. Phys.* 46, (2005), 052702.

[8] M. Corbera, J. Llibre and E. Pérez-Chavela, Equilibrium points and central configurations for the Lennard-Jones 2- and 3-body problems, *Celestial Mechanics and Dynamical Astronomy* 89(3), (2004) 235–266.

[9] J.M. Cors, G.R. Hall, G.E. Roberts, Uniqueness results for co-circular central configurations for power-law potentials, *Physica D* 280-281, (2014) 44-47.

[10] S. Craig, F. Diacu, E.A. Lacomba and E. Pérez-Chavela, On the anisotropic Manev problem, *J. Math. Phys.* 40, (1999) 1–17.

[11] F. Diacu, The planar isosceles problem for Manev’s gravitational law, *J. Math. Phys.* 34(12), (1993) 5671–5690.

[12] F. Diacu, Regularization of partial collisions in the n-body problem, *Differential Integral Equations* 5(1), (1992) 103–136.

[13] F. Diacu, Near-collision dynamics for particle systems with quasihomogeneous potentials, *J. Differential Equations* 128, (1996) 58–77.
[14] F. Diacu, On the singularities of the curved $n$-body problem, *Trans. Amer. Math. Soc.* **363** (2011) 2249–2264.

[15] F. Diacu, Polygonal homographic orbits of the curved $n$-body problem, *Trans. Amer. Math. Soc.* **364**, 5 (2012) 2783–2802.

[16] F. Diacu, Relative equilibria in the 3-dimensional curved $n$-body problem, *Memoirs Amer. Math. Soc.* **228**, 1071 (2013).

[17] F. Diacu, *Relative Equilibria of the Curved N-Body Problem*, Atlantis Studies in Dynamical Systems, vol. 1, Atlantis Press, Amsterdam, 2012.

[18] F. Diacu, The non-existence of centre-of-mass and linear-momentum integrals in the curved $n$-body problem, arXiv:1202.4739, 12 p.

[19] F. Diacu, The curved N-body problem: risks and rewards, *Math. Intelligencer* **35**, 3 (2013) 24–33.

[20] J. Delgado, F.N. Diacu, E.A. Lacomba, A. Mingarelli, V. Mioc, E. Perez, C. Stoica, The global flow of the Manev problem, *J. Math. Phys.* **37**(6), (1996) 2748–2761.

[21] F. Diacu, S. Kordlou, Rotopulsators of the curved N-body problem, *J. Differ. Equations* **255**, (2013) 2709–2750.

[22] F. Diacu, V. Mioc, and C. Stoica, Phase-space structure and regularization of Manev-type problems, Nonlinear Analysis 41, 1029–1055 (2000).

[23] F. Diacu and E. Pérez-Chavela, Homographic solutions of the curved 3-body problem, *J. Differ. Equations* **250**, (2011) 340–366.

[24] F. Diacu, E. Pérez-Chavela and M. Santoprete, Saari’s conjecture in the collinear case, *Trans. Amer. Math. Soc.* **357**, 10 (2005) 4215–4223.

[25] F. Diacu, E. Pérez-Chavela and M. Santoprete, The $n$-body problem in spaces of constant curvature, arXiv:0807.1747, 54 p.

[26] F. Diacu, E. Pérez-Chavela and M. Santoprete, The $n$-body problem in spaces of constant curvature. Part I: Relative equilibria, *J. Nonlinear Sci.* **22**, 2 (2012) 247–266.

[27] F. Diacu, E. Pérez-Chavela and M. Santoprete, The $n$-body problem in spaces of constant curvature. Part II: Singularities, *J. Nonlinear Sci.* **22**, 2 (2012) 267–275.
[28] F. Diacu, E. Pérez-Chavela, M. Santoprete, Central configurations and total collisions for quasihomogeneous $n$-body problems, *Nonlinear Analysis* **65**, (2006) 1425–1439.

[29] F. Diacu, S. Popa, All the Lagrangian relative equilibria of the curved 3-body problem have equal masses, *J. Math. Phys.* **55**, 112701 (2014).

[30] F. Diacu, B. Thorn, Rectangular orbits of the curved 4-body problem, *Proc. Amer. Math. Soc.* **143** (2015), 1583–1593.

[31] M. Gidea, J. Llibre, Symmetric planar central configurations of five bodies: Euler plus two, *Celest. Mech. Dyn. Astr.* **106** (2014), 89-107.

[32] K. G. Hadjifotinou, T. J. Kalvouridis, Numerical investigation of periodic motion in the three-dimensional ring problem of n bodies, *Int. J. Bifurcat. Chaos* **15** (2005), 2681-2688.

[33] R. Jones, Central configurations with a quasihomogeneous potential function, *J. Math. Phys.* **49**, 052901 (2008).

[34] T. J. Kalvouridis, A planar case of the n+1-body problem: The ring problem, *Astrophys. Space Sci.* **260**:309325, 1999.

[35] T. J. Kalvouridis, Zero velocity surface in the three-dimensional ring problem of n + 1 bodies, *Celest. Mech. Dyn. Astr.* **80**, 133-144, (2001).

[36] W. Killing, Die Rechnung in den nichteuklidischen Raumformen, *J. Reine Angew. Math.* **89**, (1880) 265–287.

[37] W. Killing, Die Mechanik in den nichteuklidischen Raumformen, *J. Reine Angew. Math.* **98**, (1885) 1–48.

[38] W. Killing, *Die Nicht-Euklidischen Raumformen in Analytischer Behandlung*, Teubner, Leipzig, 1885.

[39] V.V. Kozlov and A.O. Harin, Kepler’s problem in constant curvature spaces, *Celestial Mech. Dynam. Astronom.* **54**, (1992) 393–399.

[40] H. Liebmann, Die Kegelschnitte und die Planetenbewegung im nichteuklidischen Raum, *Berichte Königl. Sächsischen Gesell. Wiss., Math. Phys. Klasse* **54**, (1902) 393–423.

[41] H. Liebmann, Über die Zentralbewegung in der nichteuklidische Geometrie, *Berichte Königl. Sächsischen Gesell. Wiss., Math. Phys. Klasse* **55**, (1903) 146–153.
[42] H. Liebmann, *Nichteuklidische Geometrie*, G.J. Göschen, Leipzig, 1905; 2nd ed. 1912; 3rd ed. Walter de Gruyter, Berlin Leipzig, 1923.

[43] N.I. Lobachevsky, The new foundations of geometry with full theory of parallels [in Russian], 1835-1838, In Collected Works, V. 2, GITTL, Moscow, 1949, p. 159.

[44] J. C. Maxwell, On the stability of motions of Saturns rings. Macmillan and Cia., Cambridge, 1859.

[45] V. Mioc, M. Stavinschi, On the Schwarzschild-type polygonal (n + 1)-body problem and on the associated restricted problem, *Balt. Astron.* 7, (1998) 637-651.

[46] V. Mioc, M. Stavinschi, On Maxwell’s (n + 1)-body problem in the manev-type field and on the associated restricted problem, *Phys. Scripta* 60, (1999) 483-490.

[47] R. Moeckel, On central configurations, *Math. Z.* 205(4), (1990) 499-517.

[48] F.R. Moulton, *The Straight Line Solutions of the Problem of n Bodies*, *Ann. Math.* 12, (1910) 1–17.

[49] V. Parasciv, Central configurations and homographic solutions for the quasihomogeneous N-body problem, *J. Math. Phys.* 53, 122902 (2012).

[50] V. Parasciv, Lagrange-Pizzetti theorem for the quasihomogeneous N-body problem, *J. Math. Phys.* 53, 062901 (2012).

[51] E. Pérez-Chavela, D. Saari, A. Susin and Z. Yan, Central configurations in the charged three body problem, *Contemp. Math.* 198, 137-155 (1996).

[52] E. Pérez-Chavela and L. Vela-Arevalo, Triple collision in the quasihomogeneous collinear three-body problem, *J. Differ. Equations* 148, 186–211 (1998).

[53] D. Saari, On the role and properties of n-body central configurations, *Celestial Mech.* 21, (1980) 9-20.

[54] H. Salo, C. F. Yoder, The dynamics of coorbital satellite systems, *Astron. Astrophys.* 205 (1-2), (1988) 309-327.
[55] D.J. Scheeres, On symmetric central configurations with application to satellite motion about rings. Ph. D. Thesis. University of Michigan, 1992.

[56] D.J. Scheeres, N.X. Vinh, Linear stability of a self-gravitating ring, *Celest. Mech. Dyn. Astr.* 51:83103, 1991.

[57] E. Schering, Die Schwerkraft im Gaussischen Räume, *Nachr. Königl. Gesell. Wiss. Göttingen*, 13 July, 15 (1873), 311–321.

[58] E. Schering, Die Schwerkraft in mehrfach ausgedehnten Gaussischen und Riemannschen Räumen, *Nachr. Königl. Gesell. Wiss. Göttingen*, 26 February, 6 (1873), 149–159.

[59] S. Smale, Mathematical problems for the next century, *Mathematical Intelligencer* 20, (1998) 7–15.

[60] S. Smale, Problems on the nature of relative equilibria in celestial mechanics, *Lecture Notes in Mathematics* 197, (1971), 194–198.

[61] S. Smale, Topology and Mechanics, II, The planar n-body problem, *Invent. Math.* 11, (1970), 45–64.

[62] P. Tibboel, Polygonal homographic orbits in spaces of constant curvature, *Proc. Amer. Math. Soc.* 141, (2013), 1465–1471.

[63] P. Tibboel, Existence of a class of rotopulsators, *J. Math. Anal. Appl.* 404, (2013) 185–191.

[64] P. Tibboel, Existence of a lower bound for the distance between point masses of relative equilibria in spaces of constant curvature, *J. Math. Anal. Appl.* 416, (2014) 205–211.

[65] P. Tibboel, Existence of a lower bound for the distance between point masses of relative equilibria for generalised quasi-homogeneous n-body problems and the curved n-body problem, *J. Math. Phys.* 56, (2015) (in press).

[66] A. Wintner, *The Analytical Foundations of Celestial Mechanics*, Princeton University Press, 1941.