THEORETICAL CEPHEID PERIOD–LUMINOSITY AND PERIOD–COLOR RELATIONS IN SPITZER IRAC BANDS

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Received 2011 August 18; accepted 2011 October 20; published 2012 January 5

ABSTRACT

In this paper, the synthetic period–luminosity (P-L) relations in Spitzer’s IRAC bands, based on a series of theoretical pulsation models with varying metal and helium abundance, were investigated. Selected sets of these synthetic P-L relations were compared to the empirical IRAC band P-L relations recently determined from Galactic and Magellanic Clouds Cepheids. For the Galactic case, synthetic P-L relations from model sets with \((Y = 0.26, Z = 0.01)\), \((Y = 0.26, Z = 0.02)\), and \((Y = 0.28, Z = 0.02)\) agree with the empirical Galactic P-L relations derived from the Hubble Space Telescope parallaxes. For Magellanic Cloud Cepheids, the synthetic P-L relations from model sets with \((Y = 0.25, Z = 0.008)\) agree with both of the empirical Large Magellanic Cloud (LMC) and Small Magellanic Cloud P-L relations. Analysis of the synthetic P-L relations from all model sets suggested that the IRAC band P-L relations may not be independent of metallicity, as the P-L slopes and intercepts could be affected by the metallicity and/or helium abundance. We also derive the synthetic period–color (P-C) relations in the IRAC bands. Non-vanishing synthetic P-C relations were found for certain combinations of IRAC band filters and metallicity. However, the synthetic P-C relations disagreed with the \([3 \mu m] – [8 \mu m] P-C\) relation recently found for the Galactic Cepheids. The synthetic \([3 \mu m] – [4.5 \mu m] P-C\) slope from the \((Y = 0.25, Z = 0.008)\) model set, on the other hand, is in excellent agreement to the empirical LMC P-C counterpart, if a period range \(1.0 < \log(P) < 1.8\) is adopted.

Key words: distance scale – stars: variables: Cepheids

1. INTRODUCTION

The mid-infrared Cepheid period–luminosity (P-L, also known as the Leavitt law) relation will be important in the James Webb Space Telescope (JWST) era, as it holds the promise of deriving the Hubble constant at the ~2% level (Freedman & Madore 2010; Freedman et al. 2011). Motivated by this, Spitzer’s IRAC band \((3.6 \mu m, 4.5 \mu m, 5.8 \mu m, \text{ and } 8.0 \mu m)\) P-L relations were derived for Cepheids in our Galaxy (Marengo et al. 2010), in the Large Magellanic Cloud (LMC; Freedman et al. 2008; Ngeow & Kanbur 2008; Madore et al. 2009; Ngeow et al. 2009; Scowcroft et al. 2011), and in the Small Magellanic Cloud (SMC; Ngeow & Kanbur 2010).

The slopes of the IRAC band P-L relations are expected to be insensitive to metallicity (Freedman et al. 2008, 2011; Freedman & Madore 2010). The empirical slopes derived from a small number of Galactic Cepheids that possess parallax distances (Marengo et al. 2010), and the Magellanic Cloud Cepheids based on the OGLE-III data (see Ngeow et al. 2009; Ngeow & Kanbur 2010, for more details), are all consistent with each other. However, these slopes do not agree with the slopes derived from Galactic Cepheids that are based on the infrared surface brightness (IRSB) method (Marengo et al. 2010) and from a smaller number of LMC Cepheids (Madore et al. 2009). Therefore, the aim of this paper is to compare these empirical P-L relations to the synthetic IRAC band P-L relations based on a series of theoretical pulsating models at various metallicities and investigate the sensitivity of IRAC band P-L relations to metallicity.

A brief description of the pulsation models is given in the next section, and the synthetic P-L relations based on these models are presented in Section 3. We compared these synthetic P-L relations to their empirical counterparts and investigated the possible metallicity dependence of the synthetic IRAC band P-L relations in Sections 4 and 5, respectively. In Section 6, synthetic period–color (P-C) relations in the IRAC bands were also derived. A discussion and conclusion are presented in Section 7.

2. THE PULSATIONAL MODELS

The synthetic P-L relations adopted in this paper are based on extensive and detailed sets of nonlinear, pulsation models including a non-local, time-dependent treatment of the coupling between pulsation and convection. These models allow us to predict not only the periods and the blue boundary of the instability strip, but also the pulsation amplitudes, the detailed light and radial velocity curve morphology, and the complete topology of the strip, including the red edge (Bono et al. 1999; Fiorentino et al. 2002; Marconi et al. 2005, 2010).

For each chemical composition and mass, an evolutionary mass–luminosity (M-L) relation was adopted (see Marconi et al. 2005) and a wide range of effective temperature compositions (see, for examples, Bono et al. 1999; Fiorentino et al. 2002; Marconi et al. 2005, 2010). For each chemical composition and mass, an evolutionary mass–luminosity (M-L) relation was adopted (see Marconi et al. 2005) and a wide range of effective temperature was explored. We note that even if the effect of varying the M-L relations has been investigated for specific chemical compositions (see, for examples, Bono et al. 1999, 2000; Caputo et al. 2005) in this analysis we assume for all the chemical composition a canonical M-L relation, neglecting both mass loss and overshooting during the previous H-burning phase. This is a limitation of the adopted model sets as the above mentioned phenomena affect the P-L relations and the corresponding distance determinations (see Section 7). In total, 17 model sets with varying helium \((Y)\) and metal \((Z)\) abundance are considered in this paper, most of them are the same model sets presented in Bono et al. (2010).
The Astrophysical Journal, 745:104 (12pp), 2012 February 1

Figure 1. Example of the synthetic IRAC band P-L relations, constructed from models with \( Y = 0.25 \) and \( Z = 0.008 \).

From the resulting theoretical instability strips and the relations connecting the periods to the intrinsic stellar parameters, synthetic P-L relations have been constructed. For this purpose, we populated the predicted instability strip by adopting the procedure suggested by Kennicutt et al. (1998). In particular, \( \sim 1000 \) pulsators were uniformly distributed from the blue to the red boundary of the instability strip, with a mass law as given by \( dn/dm \sim m^{-3} \) over the mass range 5–11 \( M_\odot \) (see Caputo et al. 2000 for further details). In order to translate the pulsational properties of the investigated Cepheid models in the Spitzer IRAC bands, we have directly convolved the predicted bolometric light curves with the Spitzer filter profiles using the general integral equation (see, for example, Girardi et al. 2002)

\[
m_{S_\lambda} = -2.5 \log \left( \int_{\lambda_1}^{\lambda_2} \lambda f_\lambda S_\lambda d\lambda \right),
\]

where \( S_\lambda \) is the IRAC spectral response curve,\(^5\) \( f_\lambda \) is the stellar flux (that corresponds to model atmospheres of known \( (T_{\text{eff}}, [M/H], \log g) \)). \( f_\lambda^0 \) is the model spectrum of Vega. Concerning the model atmospheres, we have adopted the homogeneous set of updated ATLAS9 Kurucz model atmospheres and synthetic fluxes (new-ODF models).\(^6\)

### 3. THE SYNTHETIC IRAC BAND P-L RELATIONS

Figure 1 shows an example of the synthetic IRAC band P-L relations from one of the model sets described in the previous section. All the synthetic P-L relations were fitted with the form \( M_{\text{IRAC}} = a + b \times \log(P) \), restricted for pulsators within the period range \( 0.4 \leq \log(P) \leq 2.0 \) (as in Bono et al. 2010). The fitted P-L slopes \( b \) and intercepts \( a \) for each of the models are summarized in Tables 1 and 2, respectively. A majority of the model sets do not have pulsators with \( \log(P) < 0.4 \). For a few of the model sets which have a small number of short-period pulsators, the differences of the P-L slopes and intercepts between the P-L relations derived from full period range and those given in Tables 1 and 2 do not exceed 0.01. In Figure 2, the slopes of the synthetic P-L relations in various bands were compared for six of the pulsating model sets (or chemical compositions), where the linear \( BV\text{IJK} \) P-L slopes were adopted from Table 2 of Bono et al. (2010).\(^7\) As expected, the slopes of the P-L relation monotonically decrease from \( B \) to \( K \) band (see, for example, Madore & Freedman 1991; Berdnikov et al. 1996; Caputo et al. 2000; Fiorentino et al. 2002, 2007; Freedman et al. 2008; Ngeow & Kanbur 2008), and “flatten out” in the mid-infrared. However, the 4.5 \( \mu m \) P-L relations show a slight increase in their slopes when compared to the “flatter” slopes defined from 3.6 \( \mu m \) and 8.0 \( \mu m \) band P-L slopes. This slight increase of the 4.5 \( \mu m \) P-L slopes, and to some extent to the 5.8 \( \mu m \) band P-L slopes, may be explained due to the presence of CO absorption features shown in the \( \sim 4 \mu m \) to \( \sim 6 \mu m \) spectral region (for further details, see Marengo et al. 2010; Freedman et al. 2011; Scowcroft et al. 2011). In fact, the slight increase of 4.5 \( \mu m \) and 5.8 \( \mu m \) band P-L slopes was shown in all model sets presented in Table 1.

### 4. COMPARISON TO THE EMPirical P-L RELATIONS

Slopes of the current empirical IRAC band P-L relations are summarized in Table 2 of Ngeow & Kanbur (2010). These include six sets of P-L relations in our Galaxy and Magellanic Clouds. Briefly, GAL1 and GAL2 are P-L slopes derived from the “old” and “new” IRSB distances, respectively, and GAL3 are slopes derived from eight Cepheids that possess Hubble Space Telescope (HST) parallax measurements. Details concerning these Galactic P-L slopes can be found in Marengo et al. (2010). The LMC1 and LMC2 are the empirical LMC P-L slopes taken from Madore et al. (2009) and Ngeow et al. (2009), respectively, and the SMC P-L slopes were adopted from Ngeow & Kanbur (2010). As discussed in Ngeow & Kanbur (2010), these six sets of P-L slopes can be grouped to two groups characterized by steeper (\( \sim -3.46 \)) and shallower (\( \sim -3.18 \)) slopes, respectively. Both of the steeper and shallower slopes can be predicted from using \( L_\lambda = 4\pi R^2 B_\lambda(T) \) by assuming the behavior of \( B_\lambda(T) \) at long wavelengths for the IRAC bands (see Freedman et al. 2008; Neilson et al. 2010; Ngeow et al. 2010 for more details).

In addition to these six sets of empirical P-L relations obtained from random-phase observations, Scowcroft et al. (2011) have recently published LMC P-L relations based on \( \sim 80 \) Cepheids that possess mean magnitudes in 3.6 \( \mu m \) and 4.5 \( \mu m \) bands. These LMC Cepheids have been observed many times using Spitzer, with 24 evenly spaced data points per light curves, hence accurate mean magnitudes can be obtained. Their adopted P-L relations in 3.6 \( \mu m \) and 4.5 \( \mu m \) bands are denoted as LMC3 in this paper.

These empirical P-L relations can be compared to the synthetic IRAC band P-L relations from the selected model sets given in Tables 1 and 2. Bono et al. (2010) have compared the synthetic \( BV\text{IJK} \) band P-L slopes of these selected model

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\(^{5}\) Available at http://ssc.spitzer.caltech.edu/irac/

\(^{6}\) Available at http://kurucz.harvard.edu/irac/ or http://wwwuser.oat.ts.astro.it/castelli/grids.html

\(^{7}\) For consistency, we only use the linear slopes, \( b_{\text{eff}} \), from Bono et al. (2010).
Table 1
Slope of the Theoretical IRAC Band P-L Relations at Various Metallicities

| Z   | Y    | log(Z/X) | 12 + log(O/H)a | [Fe/H]b | ΔY/ΔZb | 3.6 μm | 4.5 μm | 5.8 μm | 8.0 μm |
|-----|------|----------|----------------|---------|---------|--------|--------|--------|--------|
| 0.0004 | 0.24  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.001  | 2.87  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.002  | 2.58  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.004  | 2.27  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.006  | 2.09  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.008  | 1.97  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.010  | 1.86  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.012  | 1.56  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.015  | 1.22  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.018  | 1.00  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.021  | 1.18  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.024  | 1.36  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.027  | 1.54  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.030  | 1.73  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |

Notes.

a Calculated from an online tool, http://astro.wsu.edu/models/calc/MIX.html, assuming 12 + log(O/H) = 8.66.
b ΔY/ΔZ = (Y − Yp)/Z is the relative helium enrichment ratio, where Yp = 0.23 (see Fiorentino et al. 2002, and reference therein).
c Taken from Marengo et al. (2010).

Table 2
Intercept of the Theoretical IRAC Band P-L Relations at Various Metallicities

| Z   | Y    | log(Z/X) | 12 + log(O/H)a | [Fe/H]b | ΔY/ΔZb | 3.6 μm | 4.5 μm | 5.8 μm | 8.0 μm |
|-----|------|----------|----------------|---------|---------|--------|--------|--------|--------|
| 0.0004 | 0.24  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.001  | 2.87  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.002  | 2.58  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.004  | 2.27  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.006  | 2.09  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.008  | 1.97  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.010  | 1.86  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.012  | 1.56  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.015  | 1.22  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.018  | 1.00  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.021  | 1.18  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.024  | 1.36  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.027  | 1.54  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |
| 0.030  | 1.73  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  | 0.42  |

Notes.

a Calculated from an online tool, http://astro.wsu.edu/models/calc/MIX.html, assuming 12 + log(O/H) = 8.66.
b ΔY/ΔZ = (Y − Yp)/Z is the relative helium enrichment ratio, where Yp = 0.23 (see Fiorentino et al. 2002, and reference therein).
c Taken from Marengo et al. (2010).

d to their empirical counterparts for Galactic and Magellanic Clouds Cepheids, and found they are generally in agreement. When comparing the P-L intercepts, three different values of LMC and SMC distance moduli (μLMC,SMC), respectively, were adopted (see the right panels of Figures 4 and 5). These distance moduli covered a wide range of available distance moduli in the literature for the Magellanic Clouds.

4.1. Comparison of the Galactic P-L Relations

A comparison of the synthetic and empirical Galactic P-L relations in IRAC bands is presented in Figure 3. The left panel of this figure shows that the synthetic P-L slopes from 12 + log(O/H) = 8.58 (Z = 0.01, Y = 0.26) and 12 + log(O/H) = 8.89 (Z = 0.02, Y = 0.26) model sets are in marginal agreement with the GAL2 and GAL3 P-L slopes, but not for the GAL1 P-L slopes derived from the “old” IRSB distances. However, the synthetic P-L intercepts from these two model sets agree well with the GAL3 P-L intercepts (albeit the large error bars), but disagree with the other two empirical P-L intercepts based on the IRSB methods. Marengo et al. (2010) have compared the synthetic P-L relations from 12 + log(O/H) = 8.90 (Z = 0.02, Y = 0.28) model set to the three sets of empirical Galactic P-L relations. They are also included in Figure 3. The comparison shown in this figure echoes the finding in Marengo et al. (2010) that this set of synthetic P-L...
relations agrees well with their empirical counterparts derived from the \textit{HST} parallaxes (GAL3), but not for the other two sets of empirical Galactic P-L relations.

Determination of the Galactic P-L relations is expected to improve in the near future. The Carnegie Hubble Program will observe and measure the distances to about 39 Galactic Cepheids in 3.6 $\mu$m and 4.5 $\mu$m, with 15 of them expected to have parallax measurements from \textit{Gaia} (Freedman & Madore 2010). The improved empirical Galactic P-L relations are expected to be able to discriminate the synthetic P-L relations that are best at describing the observed P-L relations.

### 4.2. Comparison of the LMC P-L Relations

In Figure 4, synthetic P-L relations from three model sets were compared to the empirical LMC P-L relations. The synthetic P-L slopes from $12 + \log(O/H) = 8.47 \ (Z = 0.008, Y = 0.25)$ model set are in good agreement with LMC2 P-L slopes from Ngeow et al. (2009). Note that the model set with $Z = 0.008$ and $Y = 0.25$ is generally adopted as a representative metallicity for LMC Cepheids. On the other hand, the synthetic P-L slopes from the $12 + \log(O/H) = 8.35 \ (Z = 0.006, Y = 0.25)$ model set also agree with LMC1 P-L slopes (from Madore et al. 2009) in 3.6 $\mu$m and 4.5 $\mu$m bands, but not in the two longer wavelength bands. Similarly, LMC3 P-L slopes from Scowcroft et al. (2011) agreed with the synthetic P-L slopes from the same model set. For the P-L intercepts, the right panels of Figure 4 show that the synthetic P-L intercepts from $12 + \log(O/H) = 8.35$ model set match with the empirical results from LMC2 if the LMC distance modulus ($\mu_{\text{LMC}}$) is adopted to be 18.50 mag. The LMC2 P-L intercepts at 5.8 $\mu$m and 8.0 $\mu$m bands also matched to the synthetic P-L intercepts from $12 + \log(O/H) = 8.47$ and $12 + \log(O/H) = 8.17$ model sets if $\mu_{\text{LMC}} \sim 18.60$ mag.

The P-L slopes adopted from Scowcroft et al. (2011) are based on Cepheids with a period range $1.0 < \log(P) < 1.8$. Table 3 of Scowcroft et al. (2011) also listed the P-L relations using Cepheids with $0.8 < \log(P) < 1.8$, at which the slopes of these P-L relations ($-3.31 \pm 0.05$ and $-3.22 \pm 0.05$ at 3.6 $\mu$m and...
Figure 4. Same as Figure 3, assuming three different values of the LMC distance modulus. These distance moduli roughly cover the available LMC distance modulus given in the literature. Note that for better visualization, data points for LMC3 have been shifted slightly in wavelength.

Figure 5. Same as Figure 3, assuming three different values of the SMC distance modulus. These distance moduli roughly cover the available SMC distance modulus given in the literature.

4.3. Comparison of the SMC P-L Relations

Synthetic P-L relations from three model sets, 12 + log(O/H) = 7.86 (Z = 0.002, Y = 0.24), 12 + log(O/H) = 8.17 (Z = 0.004, Y = 0.25), and 12 + log(O/H) = 8.47 (Z = 0.008, Y = 0.25), were compared to the empirical SMC P-L relations from Ngeow & Kanbur (2010). As can be seen from Figure 5, the synthetic P-L relations from 12 + log(O/H) = 8.47 model set agree with the empirical P-L relations, if the assumed SMC distance modulus is 19.10 mag. This result is odd because Z = 0.004 is generally adopted as representative of SMC metallicity. Though there is a spread in the metallicity of SMC Cepheids from spectroscopic measurements, the mean value is closer to Z = 0.004 than to Z = 0.008. This occurrence is
Figure 6. Synthetic P-L relations as a function of 12 + log(O/H), separated by the helium abundance (Y). The left and right panels are for the P-L slopes and intercepts, respectively.

Figure 7. Synthetic P-L relations as a function of 12 + log(O/H), separated according to the ΔY/ΔZ values. Note that the open symbols are for ΔY/ΔZ > 3, and filled symbols are for ΔY/ΔZ between 1 and 3 (inclusive).

due to the fact that current nonlinear pulsation models predict a significant steepening of P-L slopes when changing the metallicity from Z = 0.008 to Z = 0.004 (with a significant reduction of the dependence at still smaller metal contents) whereas the empirical relations adopted in this paper tend to suggest almost the same slopes for LMC and SMC. For 12 + log(O/H) = 8.17 model set, the synthetic P-L slopes agree with the empirical P-L relations at 5.8 μm and 8.0 μm bands, though the 3.6 μm and 4.5 μm P-L slopes are slightly off with respect to their empirical counterparts.

5. DEPENDENCY OF THEORETICAL P-L RELATIONS ON METALLICITY

As mentioned in the Introduction, the IRAC band P-L relations are expected to be insensitive to metallicity. However, based on the residual analysis from multi-band empirical P-L relations for a number of LMC Cepheids with [Fe/H] measurements, Freedman & Madore (2011) suggested the mid-infrared P-L relation could be mildly dependent on metallicity. Therefore, possible metallicity dependence of the IRAC band P-L relations is investigated in this section using the synthetic P-L relations presented in Tables 1 and 2. Figures 6 and 7 present the synthetic P-L slopes (left panels) and intercepts (right panels) as a function of 12 + log(O/H), suggesting the IRAC band P-L relations may be sensitive to metallicity and/or helium abundance of the Cepheids.

Well-known statistical tests can be used to test the possible metallicity dependence of the synthetic P-L relations as shown in Figures 6 and 7. We considered a null hypothesis that the IRAC band P-L relations are independent of metallicity, which can be represented by a constant regression model in the form of A = a₀, where A stands for either the P-L slopes or intercepts. An alternative hypothesis is that there is a linear dependence of

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9 Plots for the synthetic P-L relations as a function of log(Z/X) and [Fe/H] measurements, Freedman & Madore (2011) suggested the mid-infrared P-L relation could be mildly dependent on metallicity. Therefore, possible metallicity dependence of the IRAC band P-L relations is investigated in this section using the synthetic P-L relations presented in Tables 1 and 2. Figures 6 and 7 present the synthetic P-L slopes (left panels) and intercepts (right panels) as a function of 12 + log(O/H), suggesting the IRAC band P-L relations may be sensitive to metallicity and/or helium abundance of the Cepheids.

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metallicity on P-L relations represented by a linear regression model, \( A = a_0 + a_1B \) where \( B = 12 + \log(O/H) \) or \( B = [\text{Fe/H}] \). Graphic representations of these regression models are presented in Figure 8. It is worth pointing out that these regression models are used to test the dependence of metallicity on synthetic P-L relations and do not represent the actual linear dependence of metallicity on P-L relations, since the metallicity effect, if present, will not be a simple linear relation (as evident in Figures 6 and 7). It is well known in statistical literature (for example, see Kutner et al. 2005) that the \( F \)-test can be applied to test if an additional parameter (in our case, \( a_1 \)) is needed in the regression model. By adopting \( \alpha = 0.05 \), the null hypothesis can be rejected if \( F > 4.54 \). The \( F \)-test results are summarized in Table 3, showing that the null hypothesis can be rejected in most cases, except for 4.5 \( \mu m \) and 5.8 \( \mu m \) band P-L intercepts. These suggested that the synthetic P-L relations may not be independent of metallicity.

6. THE SYNTHETIC IRAC BAND P-C RELATIONS

Pulsators from various model sets, as described in Section 2, can also be used to construct the synthetic P-C relations in the IRAC bands. For brevity, colors from 3.6 \( \mu m \) and 4.5 \( \mu m \) bands are denoted as [3.6]–[4.5], and so on. Two examples of the synthetic P-C relations are presented in Figure 9. Results of the synthetic P-C relations are summarized in Table 4. As in the case of the synthetic P-L relations, pulsators with \( 0.4 \leq \log(P) \leq 2.0 \) were used to fit the P-C relations. The P-C slopes and intercepts do not deviate by more than 0.006 if all the pulsators were included. The synthetic P-C slopes and intercepts were plotted as a function of 12 + \( \log(O/H) \).
in Figures 10 and 11, respectively. From these figures, it is clear that P-C relations exist for certain combinations of IRAC band filters and metallicity, and not all of the synthetic P-C relations are independent of metallicity. The P-C relations that are independent or insensitive to metallicity are the $[3.6]-[8.0]$ and $[4.5]-[5.8]$ P-C relations, especially for those with $12 + \log(O/H) < 8.9$.

### 6.1. Comparison to Empirical P-C Relations

Marengo et al. (2010) found a significant $[3.6]-[8.0]$ P-C relation for the Galactic Cepheids, with the expression of $[3.6]-[8.0] = 0.039(\pm0.008) \log(P) - 0.058(\pm0.014)$. This empirical non-zero P-C slope is in contradiction with the synthetic P-C slopes given in Table 4, at which the synthetic $[3.6]-[8.0]$ P-C slopes are close to zero. Disagreement was also found for the P-C intercepts. In Figure 12, P-C relations for 26 fundamental model Galactic Cepheids are compared to the synthetic P-C relations from selected model sets. IRAC band photometry for these Galactic Cepheids was adopted from Marengo et al. (2010). This figure shows that only the $[3.6]-[4.5]$ and $[5.8]-[8.0]$ synthetic P-C relations barely agreed with the observed P-C relations. Disagreements between empirical and synthetic P-C relations may be due to the small number of Cepheids in the sample, large photometric errors, colors for all Cepheids are not be measured at mean light, presence of circumstellar envelopes especially for longer period Cepheids, uncertainties in the adopted model atmospheres, or any combinations of these.

In addition to Galactic Cepheids, Scowcroft et al. (2011) also reported the finding of a $[3.6]-[4.5]$ P-C relation based on LMC Cepheids, within the period range $1.0 < \log(P) < 1.8$, that possess multi-epoch observations. The left panel of Figure 13 compares the P-C relation for these Cepheids and the three synthetic P-C relations given in Table 4, where the $[3.6]-[4.5]$ colors are obtained from the $3.6\mu$m and $4.5\mu$m band mean magnitudes taken from Scowcroft et al. (2011). It is clear that these synthetic P-C relations do not agree with the empirical P-C relation of $[3.6]-[4.5] = -0.087(\pm0.012) \log(P) + 0.092(\pm0.013)$ from Scowcroft et al. (2011). However, as in Section 4.2, this empirical P-C relation should also be compared to the synthetic P-C relations based on pulsators within the period range $1.0 < \log(P) < 1.8$. These synthetic P-C relations for the three model sets are summarized in Table 5 and compared to the empirical P-C relation in the right panel of Figure 13. Table 5 shows that the synthetic P-C slope from $(Z = 0.008, Y = 0.25)$ model set is in excellent agreement with the empirical P-C slope, although the synthetic P-C intercept is $\sim2\sigma$ smaller than its empirical counterpart. On the other hand, the synthetic P-C relation from $(Z = 0.006, Y = 0.25)$ model set is consistent with the empirical P-C relation.

### 6.2. The Wesenheit Function

The existence of the P-C relations suggested the Wesenheit function can be formulated in the IRAC bands. In optical bands,
Figure 11. Same as Figure 10, but for the intercepts of the synthetic P-C relations.

Table 4  
Synthetic IRAC Band P-C Relations at Various Metallicities$^a$

| $Z$   | $Y$  | [3.6]–[4.5] | [3.6]–[5.8] | [3.6]–[8.0] | [4.5]–[5.8] | [4.5]–[8.0] | [5.8]–[8.0] |
|-------|------|-------------|-------------|-------------|-------------|-------------|-------------|
| 0.0004| 0.24 | −0.005      | 0.002       | 0.008       | 0.007       | 0.012       | 0.006       |
| 0.001 | 0.24 | −0.021      | −0.008      | 0.008       | 0.013       | 0.029       | 0.016       |
| 0.002 | 0.24 | −0.024      | −0.011      | 0.008       | 0.013       | 0.032       | 0.019       |
| 0.004 | 0.25 | −0.039      | −0.023      | 0.006       | 0.016       | 0.045       | 0.029       |
| 0.006 | 0.25 | −0.061      | −0.039      | 0.006       | 0.023       | 0.067       | 0.045       |
| 0.008 | 0.25 | −0.057      | −0.038      | 0.004       | 0.019       | 0.061       | 0.042       |
| 0.01  | 0.26 | −0.062      | −0.043      | 0.003       | 0.019       | 0.065       | 0.045       |
| 0.02  | 0.25 | −0.072      | −0.053      | 0.001       | 0.019       | 0.073       | 0.054       |
| 0.02  | 0.26 | −0.065      | −0.047      | 0.003       | 0.018       | 0.067       | 0.050       |
| 0.02  | 0.28 | −0.096      | −0.068      | −0.01       | 0.034       | 0.068       | 0.057       |
| 0.03  | 0.31 | −0.079      | −0.058      | 0.001       | 0.021       | 0.081       | 0.059       |
| 0.03  | 0.31 | −0.086      | −0.061      | −0.012      | 0.025       | 0.074       | 0.049       |
| 0.03  | 0.335| −0.088     | −0.063       | −0.012     | 0.025       | 0.075       | 0.051       |
| 0.04  | 0.25 | −0.075      | −0.055      | −0.007     | 0.020       | 0.068       | 0.048       |
| 0.04  | 0.29 | −0.078      | −0.057      | −0.010     | 0.022       | 0.068       | 0.047       |
| 0.04  | 0.33 | −0.078      | −0.057      | −0.011     | 0.021       | 0.067       | 0.045       |

| $Z$   | $Y$  | [3.6]–[4.5] | [3.6]–[5.8] | [3.6]–[8.0] | [4.5]–[5.8] | [4.5]–[8.0] | [5.8]–[8.0] |
|-------|------|-------------|-------------|-------------|-------------|-------------|-------------|
| 0.0004| 0.24 | 0.008       | 0.008       | 0.014       | −0.009      | 0.006       | 0.006       |
| 0.001 | 0.24 | 0.020       | 0.015       | 0.012       | −0.005      | −0.008      | −0.003      |
| 0.002 | 0.24 | 0.015       | 0.012       | 0.011       | −0.003      | −0.004      | −0.001      |
| 0.004 | 0.25 | 0.022       | 0.019       | 0.012       | −0.003      | −0.010      | −0.007      |
| 0.006 | 0.25 | 0.034       | 0.027       | 0.011       | −0.006      | −0.023      | −0.016      |
| 0.008 | 0.25 | 0.022       | 0.021       | 0.012       | −0.001      | −0.011      | −0.009      |
| 0.01  | 0.26 | 0.017       | 0.019       | 0.013       | 0.002       | −0.004      | −0.006      |
| 0.02  | 0.25 | 0.001       | 0.008       | 0.013       | 0.007       | 0.012       | 0.005       |
| 0.02  | 0.26 | 0.002       | 0.008       | 0.011       | 0.007       | 0.010       | 0.003       |
| 0.02  | 0.28 | 0.01b       | 0.03b       | 0.02b       | 0.02b       | 0.01b       | −0.01b      |
| 0.02  | 0.31 | 0.019       | 0.022       | 0.012       | 0.002       | −0.007      | −0.009      |
| 0.03  | 0.275| −0.011      | −0.001      | 0.017       | 0.010       | 0.028       | 0.018       |
| 0.03  | 0.31 | 0.002       | 0.006       | 0.022       | 0.004       | 0.021       | 0.016       |
| 0.03  | 0.335| 0.006       | 0.010       | 0.023       | 0.004       | 0.017       | 0.013       |
| 0.04  | 0.25 | −0.020      | −0.009      | 0.018       | 0.011       | 0.038       | 0.027       |
| 0.04  | 0.29 | −0.015      | −0.006      | 0.020       | 0.010       | 0.035       | 0.026       |
| 0.04  | 0.33 | −0.012      | −0.003      | 0.022       | 0.009       | 0.035       | 0.026       |

Notes.
$^a$ Errors of the P-C slopes and intercepts are ignored as they are generally less than 0.003.
$^b$ Derived from the P-L relations.
dispersion of the Wesenheit function is \( \sim 2 \) to \( \sim 3 \) times smaller than the optical P-L relations (for example, see Fouqué et al. 2007; Soszynski et al. 2008; Ngeow et al. 2009, and references therein). The dispersions of the synthetic Wesenheit function, in the form of \( W = m_{3.6\mu m} - 3.40 \times ([3.6] - [4.5]) \), from all model sets are on average about \( \sim 1.5 \) larger than the dispersions from the synthetic 3.6\( \mu m \) band P-L relations. Also, it is not worth deriving the IRAC band Wesenheit function because the Wesenheit function is reddening free by definition. On the other hand, extinction can almost be ignored for the IRAC band.
P-L relations (Freedman et al., 2008, 2011; Freedman & Madore, 2010; Ngeow et al., 2009; Marengo et al., 2010). Therefore, there is no net gain in using the Wesenheit function in the IRAC bands.

7. DISCUSSION AND CONCLUSION

Synthetic IRAC band P-L relations based on a series of pulsation models were investigated in this paper. These synthetic P-L relations were compared to their empirical counterparts, and the possible metallicity dependency of the IRAC band P-L relations was examined. For the former part, selected sets of synthetic IRAC band P-L relations show agreement to the empirical P-L relations derived from Galactic and Magellanic Cloud Cepheids. The $BVIJK$ synthetic P-L relations for these selected model sets also agreed with their optical and near-infrared counterparts, as presented in Bono et al. (2010), for the Galactic and Magellanic Clouds P-L relations. For the metallicity dependency of IRAC band P-L relations, plots for all of the synthetic P-L relations as a function of metallicity revealed that the IRAC band P-L relations may not be independent of metallicity. This result is also supported by a statistical $F$-test. On the other hand, current empirical IRAC band P-L relations based on three galaxies, either the P-L relations in “steep” or “shallow” groups (see Table 2 of Ngeow & Kanbur, 2010), suggested the IRAC band P-L relations may not depend on metallicity, or have a weak dependency on metallicity. It is clear that more empirical determinations of the IRAC band P-L relations are needed, especially for galaxies with low metallicity, from future observations with JWST.

Figure 14 shows the intercepts for LMC and SMC empirical P-L relations relative to the two adopted synthetic P-L relations, which is equivalent to deriving the distance moduli to the Magellanic Clouds using the synthetic P-L relations. The resultant distance moduli for LMC and SMC are higher than the values commonly adopted in the literature (e.g., 18.6 mag and 19.6 mag for LMC and SMC, respectively). This is due to the already mentioned limitation of the adoption of a canonical M-L relation. However, previous theoretical computations of current nonlinear convective models (see, e.g., Bono et al., 2002, 2008; Caputo et al., 2002, and references therein) have shown that by relying on non-canonical models based on an M-L relation brighter than the canonical one by 0.25 dex the inferred distance moduli are shorter than the values obtained in the canonical scenario by 0.15 to 0.2 mag depending on the filters. This implies that adopting non-canonical relations in the comparison with empirical Spitzer data we would have obtained shorter distance moduli by 0.15 to 0.2 mag for each selected chemical composition, in better agreement with most recent and adopted values in the literature.

Finally, synthetic P-C relations in the IRAC band were also derived and compared to their Galactic and LMC empirical counterparts. In general, disagreements were found between the synthetic and empirical P-C relations. However, the synthetic P-C relations are in agreement with the empirical LMC [3.6]–[4.5] P-C relation if a period range $1.0 < \log(P) < 1.8$ is adopted when constructing the synthetic P-C relations. Observations of a large number of Cepheids with JWST may help in resolving the discrepancy of the IRAC band P-C relations.

The authors thank the referee for constructive comments to improve this manuscript. C.C.N. thanks the funding from National Science Council (of Taiwan) under the contract NSC 98-2112-M-008-013-MY3. This work is based (in part) on observations made with the Spitzer Space Telescope, which is operated by the Jet Propulsion Laboratory, California Institute of Technology under a contract with NASA.

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