ABSTRACT: As we know, oxygen usually forms two single bonds with other atoms, whereas in this work, we reported the neutral molecules with square planar O(B)4-type tetracoordinate oxygen substructures. The difficulty to achieve a square planar O(B)4-type arrangement is not only to overcome the strain from the planar to tetrahedral configuration but also to constrain it in a right system with the proper symmetry. Several neutral molecules with square planar O(B)4-type arrangements have been studied using the DFT method at the B3LYP/6-311++G(3df,3pd) level of theory. The computational results show that they are all real minima on potential energy hypersurfaces. Thus, it has been shown theoretically that the square planar O(B)4-type arrangement is achieved. The molecular orbitals among the square planar O(B)4-type substructure at $D_{2d}$ symmetry have been suggested. Using the molecular orbitals, it has been explained that the oxygen 2p$_z$-$\pi$ electrons in the square planar O(B)4-type arrangement are no longer the lone-pair electrons.

INTRODUCTION

Oxygen occupies a central position in inorganic chemistry. It has four valence orbitals and six valence electrons; therefore, it usually forms two single bonds with other atoms. When an oxygen atom bonds to four or more atoms in the same plane, the oxygen is called a planar hypercoordinate oxygen (phO). Species with planar tetracoordinate oxygen (ptO) arrangements are the most studies among phOs. In the past three decades, species with planar hypercoordinate oxygen atoms have been made notable progress. In 1991, the species Al$_4$O with planar tetracoordinate oxygen and $D_{4h}$ symmetry was found by Schleyer and Boldyrev. The hydrometal complexes Cu$_4$H$_4$O$_2^+$ with ptO(Cu)$_4$ and Ni$_4$H$_4$O$_2^+$ with ptO(Ni)$_4$ substructures were studied by Li and co-workers in 2004. A planar pentacoordinate oxygen, the hydrometal complexes Cu$_5$H$_5$O$_2^+$ with a ptO(Cu)$_5$ substructure, also has been reported in 2005. A perfectly square-planar tetracoordinate oxygen atom in a tetracopper cluster-based coordination polymer was experimentally detected by Zhang and co-workers in 2011.

Most of the successful examples of ptOs are used metal-based groups as ligands previously. The achievement of ptO-used nonmetal-based groups as ligands, where the central oxygen is surrounded by four nonmetal-based groups, is more challenging. Until now, to our knowledge, no structure with a ptO(X)$_4$ (X = nonmetal atom) substructure has been reported. Among all ptO(X)$_4$ (X = nonmetal atom) substructures, the configuration of the highest symmetry is the square arrangement. It is intriguing to ask whether a square ptO(X)$_4$ (X = nonmetal atom) arrangement could exist or not? In order to answer the question, we will study neutral molecules with square ptO(B)$_4$ arrangements here.

As to design species with square ptO(B)$_4$ arrangement, we used three strategies in this work. The first is that using electron-deficient boron as donor atoms to compensate electron-rich oxygen. Because oxygen has six valence electrons and four valence orbitals, when it bonds with four atoms, the better approach is bonding with an electron-deficient element, such as boron. Thus, we used four boron-based groups as ligands to achieve the square ptO(B)$_4$ substructure.

The second is that constraining a ptO(B)$_4$ arrangement in a more buttress cage-like molecule to overcome the strain of the planar-tetracoordinate oxygen to transfer to a tetrahedral substructure, such as oxygen in the tetrahedral substructure of...
Zn\textsubscript{4}O in MOF-5\textsuperscript{a,b} and Be\textsubscript{4}O in “basic beryllium acetate”\textsuperscript{.5b}
Radom and co-workers\textsuperscript{6,7} used the “mechanical” strategy to design a family of buttress alkaplane molecules, which constricted the planar tetracoordinate carbon (ptC) arrangements in hydrocarbon cages. In 2002, Wang and Schleyer\textsuperscript{8} stated that “although approaching a planar C(C)\textsubscript{4}-type ptC more closely than ever before, such “mechanical” designs without “electronic” assistance must struggle hard to overcome the enormous strain of a ptC with a p-π lone pair HOMO.” Therefore, the candidates of cage-like molecules with more buttress than those of the alkaplane molecules designed by Radom and co-workers are required.

The third is selecting a more buttress cage-like molecule with the proper symmetry. To achieve a square ptO(B)\textsubscript{4} arrangement, except to overcome the enormous strain from the planar to tetrahedral, another problem is that the more buttress cage-like molecule should have a proper symmetry to make the ptO(B)\textsubscript{4} arrangement squarely.

Because the buttresses of the cage-like molecules are not enough, the singlet neutral molecule containing the square ptO(B)\textsubscript{4} substructure could not be achieved with the structures of the alkaplane designed by Radom and co-workers\textsuperscript{6,7} and the boraplane designed by Wang and Schleyer.\textsuperscript{9} We considered another cage-like molecule, isopagodane C\textsubscript{20}H\textsubscript{20}.\textsuperscript{9} Because there are two C=C−C=C bonds in isopagodane (Figure 1), the structure of isopagodane is different from that of the alkaplane designed by Radom and co-workers\textsuperscript{6,7} but more rigid than the structures described above. Based on the isopagodane-like structure, we designed the singlet neutral molecules 1–6 containing the square ptO(B)\textsubscript{4} substructures (Figure 1). Using four boron atoms and four BH groups to replace the four carbon atoms at the center and the four CH\textsubscript{2} groups in isopagodane, 1 was obtained. Based on the structure of 1, using four AlH, Be, CH\textsubscript{3}, NH, and O groups to replace the four BH groups in 1, 2–6 have been obtained.

In order to verify whether 1–6 are candidates with successfully square ptO(B)\textsubscript{4} substructures or not, they have been studied using the DFT method.

### RESULTS AND DISCUSSION

**Geometries and Stability.** The optimized geometries of 1–6 are all converged with D\textsubscript{2h} point groups at the B3LYP/6-311++G(3df,3pd) level of theory. Based on D\textsubscript{2h} point groups, all bond angles of ∠BOB are equal to 90°. Therefore, all ptO(B)\textsubscript{4} arrangements in 1–6 are all perfectly square planar tetracoordinate oxygen arrangements. The smallest vibrational frequencies of 1, 2, and 3 are 170.6, 127.7, and 180.8 cm\textsuperscript{-1}, respectively, without any imaginary vibrational frequency, which confirms that 1–3 correspond to true minima on the potential energy hypersurfaces (PES). In contrast, each of 4–6 has two imaginary vibrational frequencies, verifying that 4–6 could not be stable molecules.

There are two ranges of vibrational frequencies of 1, 170–1316 cm\textsuperscript{-1}, the vibrations of the atoms on the skeleton of cage, and 2632–3127 cm\textsuperscript{-1}, the vibrations of the B−H and C−H bonds. As for 2, the vibrational frequencies of the atoms on the skeleton of cage are among 127–1312 cm\textsuperscript{-1} and the vibrational frequencies of the Al−H and C−H bonds are in the 1922–3115 cm\textsuperscript{-1} range. The vibrational frequencies of the atoms on the skeleton of cage and the C−H bonds of 3 are in the 180–1303 and 3062–3147 cm\textsuperscript{-1} ranges, respectively.
The optimized structures of 1–6 with key parameters are shown in Figure 1. It can be seen that the lengths of O–B bonds in square ptO(B)₄ arrangements for 1, 2, 3, 4, 5, and 6 are 1.5390, 1.5624, 1.5634, 1.4984, 1.5008, and 1.4991 Å, respectively, which are all longer than those of single O–B in boric acid B(OH)₃ (1.376 Å). From the bond lengths, it could be seen that the O–B bonds in 4–6 are stronger than those in 1–3. Because the stronger the O–B bonds, the greater the energy difference between the tetrahedral and planar O(B)₄ arrangements as the ptC, the planar geometries of 1–3 may be more stabilized or the tetrahedral more destabilized than those of 4–6.

In the structures of 1–3, the lengths of B₂–C₃ (see the labels in Figure 1, similarly hereinafter) bonds, 1.5758, 1.5647, and 1.5626 Å for 1, 2, and 3, respectively, are all shorter than those of the B–C bond in trimethylborane (CH₃)₃B (1.578 Å), which makes the structures of 1–3 more buttress, achieving the square ptO(B)₄ arrangements successfully; while the lengths of B₂–C₃ bonds, 1.5971, 1.6050, and 1.6057 Å for 4, 5, and 6, respectively, are all longer than that of the B–C bond in (CH₃)₃B (1.578 Å). Thus, the “mechanical” strength or rigidity around the square ptO(B)₄ arrangements of 4–6 is all lower than that of 1–3. Therefore, the square planar substructures of O(B)₄ in 4–6 are not the stable, 4–6 are the saddle points on the PES, not the stable molecules. According to the computational results, the differences of energies between the saddle points and the stable structures of 4–6 are 13.0, 60.3, and 67.9 kJ·mol⁻¹, respectively.

As we know, a small HOMO–LUMO gap has long been recognized as being correlated with reactivity and structural instability, while a large gap is associated with kinetic and structural stability. The HOMO–LUMO gaps of 1–3 are 3.91, 3.16, and 2.64 eV at the B3LYP/6-311++G(3df,3pd) level of theory, respectively, and they are all larger than that of the boric acid B(OH)₃ (1.376 Å). Hence, 1–3 are all stable than the boric acid CₓBₙHₘ, Therefore, our work could be used to achieve the planar-tetracoordinate oxygen.

From the results described above, the greater “mechanical” strength or rigidity and the lower destabilized planar geometries of the square ptO(B)₄ arrangements are the key reasons to achieve the square ptO(B)₄ arrangements.

**Bonding Interactions of the Square ptO(B)₄ Arrangement.** The molecular orbitals of planar methane have been proposed using the molecular symmetry and group theory by Hoffmann, Alder, and Wilcox in 1970.11 We could use the method to discuss the nature of the bonding interactions in the square ptO(B)₄ arrangement. It could be roughly considered that every B₂ atom uses two sp hybrids, formed by 2s and 2pₓ, and two electrons to bond with two C₃ atoms and leave 2pᵧ and 2pₐ orbitals and one electron to bond with the central O₁ atom.

The representations of the basis sets formed by the eight orbitals, 2pₓ and 2pᵧ of four B₂ atoms, could be reduced to Γ₁₂₃ = a₁ + a₂ + b₁ + b₂ + 2e at the D₂h symmetry of the whole molecule. The a₁, a₂, b₁, and b₂ are all one-dimensional, and the e is two-dimensional irreps. According to the a₁, a₂, b₁, b₂, and 2e irreps, the ligand group orbitals formed by the 2pₓ and 2pᵧ of four B₂ atoms could be derived out as the symmetry-adapted linear combinations to interact with central O₁ atomic valence orbitals of the same symmetry, and form the molecular orbitals (MOs). The central O₁ atomic valence orbitals belong to a₁(2s), b₁(2pₓ), and e(2pᵧ and 2pₐ) symmetries. Thus, the diagram of various orbitals in the ptO(B)₄ substructure have been shown in Figure 2. Because every B₂ atom uses two electrons to form two B₂–C₃ bonds and leaves only one electron to bond with the central O₁ atom and there are no formal charges on four B₂ and O₁ atoms, a total of ten electrons occupy the five MOs according to the Aufbau principle. According to the electronegativities of boron and oxygen, the MOs of 1a₁, 1e, and 1b₁ possess predominantly oxygen 2s, 2pₓ, and 2pᵧ characters, respectively, and 2b₂ MO possesses predominantly boron 2pᵧ and 2pₓ characters (see HOMO(b₂) in Figure 3). Thus, the six electrons of oxygen occupy 1a₁, 1e, and 1b₁ MOs predominantly, and the four electrons of the four perimeter boron atoms occupy 1e and 2b₂ MOs, respectively.

Because the central O₁ atomic 2pₓ orbital has the b₂ irrep, it could interact with the ligand group orbitals of the b₂ symmetry to form the bonding 1b₂ and 2b₂ MOs; thus, the 2pₓ-π electrons of oxygen in the square ptO(B)₄ arrangement are no longer the lone-pair electrons (see the 1b₂ and 2b₂ MOs in Figure 2). However, oxygen is an electron-rich atom, and using the electron-deficient element, such as boron, as donor
atom may delocalize the electron density of the central O1 atom, which could decrease the energy of the bonding MOs containing the $2p_z$ electrons of the central O1 atom and stabilize the ptO(B)$_4$ substructure. Thus, the “electronic” approach, stabilized by π acceptor substituents or by aromatic delocalization and benefited from σ donation by electropositive groups, for designing the ptC compounds is also suitable for the ptO substructure.

Figure 3 shows the frontier orbitals of 1 at the B3LYP/6-311+G(3df,3pd) level of theory for the view of comparing. It could be seen that the HOMO and LUMO of Figure 2.

Because the oxygen atom form four O–B bonds at the same plane, the Wiberg bond indexes of O–B bonds in ptO(B)$_4$ are low. According to the $\rho(r)_s$, 0.1084–0.1240 a.u., in Table 2, the strength of the O–B bonds are similar to BrF and IF.

Although the distances between B2 and B2′, 2.11–2.21 Å for 1–6, are much longer than the length of the B–B bond (1.590 Å for B$_2$), some bonding interactions among perimeter four boron atoms in the square ptO(B)$_4$ still remain according to the MOs given in Figure 2. This can be seen from the Wiberg bond indexes of B2...B2′ and B2...B2″ in Table 1, and the bonding interactions among perimeter four boron atoms are much stronger than expected. The unexpected strong bonding interactions among perimeter four boron atoms bring the square ptO(B)$_4$ substructure more rigidity than expected, which is another reason to achieve the square ptO(B)$_4$ arrangement.

**Characterization of Bonding Interactions.** In order to understand the bonding interactions, we used the critical points (CPs) and their properties of the quantum theory of atoms in molecules (QTAM) to quantify the bonding interactions of atoms in 1–6, which were performed using the program Multiwfn. In QTAM, the critical points (CPs) are the points with a vanishing gradient of the electron density $\rho(r)$, that is, $\nabla \rho(r) = 0$ and could be divided into nuclear, bond, ring, and cage CPs, depending on the nature of the extremum of $\rho(r)$. When $n$, $b$, $r$, and $c$ used to represent numbers of nuclear, bond, ring, and cage CPs, they are governed by the Poincaré–Hopf relationship as $n = b + r - c = 1$. All CPs in 1–6 have been searched out using the Multiwfn program. They are all satisfied the Poincaré–Hopf relationship. All CPs and the bond paths of 1–6 have been shown in Figure 4.

The properties of CPs, electron densities $\rho(r)_s$, Laplacian of electron densities $\nabla^2 \rho(r)_s$, the densities of kinetic energy $G(r)_s$, potential energy $V(r)_s$, total electronic energy densities

| properties | 1 | 2 | 3 | 4 | 5 | 6 |
|------------|---|---|---|---|---|---|
| ES | $^1A_1$ | $^1A_1$ | $^1A_1$ | $^1A_1$ | $^1A_1$ | $^1A_1$ |
| PG | $D_{2d}$ | $D_{2d}$ | $D_{2d}$ | $D_{2d}$ | $D_{2d}$ | $D_{2d}$ |
| freq | 170.6 | 127.7 | 180.8 | 121.2i | 215.6i | 219.2i |
| WBL1O–B | 0.3321 | 0.3414 | 0.3510 | 0.3589 | 0.3641 | 0.3697 |
| WBL1C–C | 0.8743 | 0.9652 | 0.9705 | 0.8392 | 0.8602 | 0.8689 |
| WBL1C–X | 0.9652 | 0.9821 | 0.9681 | 0.9736 | 0.9739 | 0.9704 |
| WBL1C–X′ | 0.8654 | 0.5146 | 0.5915 | 0.9976 | 0.9540 | 0.8818 |
| WBL1C–X″ | 0.9775 | 0.9967 | 0.9518 | 0.9739 | 0.9730 | 0.9762 |
| WBL1C–B′ | 0.1084 | 0.1849 | 0.1791 | 0.2582 | 0.2559 | 0.2513 |
| WBL1C–B′′ | 0.0704 | 0.0635 | 0.0580 | 0.1545 | 0.1519 | 0.1544 |
| tWBL1O | 1.5880 | 1.6404 | 1.6659 | 1.6748 | 1.6949 | 1.7094 |
| tWBL1C | 2.9703 | 3.1151 | 3.1134 | 2.9739 | 3.0230 | 3.0256 |
| tWBL1C′ | 3.7320 | 3.5705 | 3.5489 | 3.8342 | 3.8300 | 3.7673 |
| tWBL1C″ | 3.9979 | 4.0103 | 3.9984 | 3.9925 | 3.9901 | 3.9844 |
| tWBL1X′ | 3.0711 | 2.1932 | 1.4516 | 3.9421 | 2.9703 | 1.9890 |
| Q1O–B1O | −1.1421 | −1.1052 | −1.0767 | −1.1489 | −1.1267 | −1.1166 |
| Q1O–B2O | 0.9731 | 0.7933 | 0.7905 | 0.9584 | 0.8896 | 0.8799 |
| Q1O–B1C3 | −0.8974 | −1.0911 | −1.0905 | −0.6015 | −0.3935 | −0.2227 |
| Q1O–B1C′ | −0.1633 | −0.2007 | −0.1994 | −0.1758 | −0.1852 | −0.2238 |
| Q1O–B1X′ | 0.5949 | 1.5207 | 1.1491 | −0.3279 | −0.6352 | −0.5828 |
| Q1O–2X′ | 1.5730 | 1.5694 | 1.5677 | 1.5671 | 1.5691 | 1.5701 |
| Q1O–2X′′ | 1.9551 | 1.9508 | 1.9529 | 1.9563 | 1.9570 | 1.9571 |
| Q1O–2X″ | 1.7872 | 1.7739 | 1.7619 | 1.7781 | 1.7682 | 1.7627 |
| Q1O–2X″′ | 1.7872 | 1.7739 | 1.7619 | 1.7781 | 1.7682 | 1.7627 |

*Labels of atoms of 1–6 given in Figure 1.*
$E_b(r_c) = G(r_c) + V(r_c)$, and the ellipticities of all bonds in the cage skeletons of $1$–$6$ at the B3LYP/6-311++G(3df,3pd) level of theory are listed in Table 2.

It can be seen that all B$_2$–C$_3$, C$_3$–C$_4$, C$_4$–C$_4'$, and C$_3$–X$_5$ (X$_5$ = B$_5$, C$_5$, N$_5$, O$_5$) bonds in $1$–$6$ are typical covalent bonds, according to the classifications,$^{16,17}$ but C$_3$–Al$_5$ in $2$ and C$_3$–Be$_5$ bonds in $3$ could be regarded as charge-shift bonds. As for O$_1$–B$_2$ bonds in the square ptO(B)$_4$ substructures of $1$–$6$, although the $\rho(r_c)s$, $0.1084$–$0.1240$ a.u., are all relatively large, the $V(r_c)s$ are all slightly greater than the corresponding $G(r_c)s$, so the values of $E_b(r_c)s$ are all negative and low, and the $V^2\rho(r_c)s$ all with positive values. Like BrF and IF,$^{17}$ these bonds have features of intermediate interaction and could be considered as partially covalent and partially electrostatic, or strongly polar bonds. The ellipticities in Table 2 also reveal that O$_1$–B$_2$ bonds have large $\pi$ bonding interactions.

Figure 4 shows that there are some weak interactions between O$_1$ and C$_4$ in $1$–$3$, but $4$–$6$ do not have these weak interactions (see $4$–$6$ in Figure 4). Table 2 gives the CPs' properties of these weak interactions. The $\rho(r_c)s$ ($0.0305$–$0.0312$ a.u.), the $E_b(r_c)s$ ($0.0027$–$0.0033$ a.u.), and the $V^2\rho(r_c)s$ ($0.1509$–$0.1606$ a.u.) are all in accordance with the typical weak interactions.$^{19}$ These weak interactions are the one of reasons to achieve the square ptO(B)$_4$ arrangement.

**CONCLUSIONS**

The singlet neutral molecules $1$–$6$ with the square planar-tetracoordinate oxygen surrounded by four boron-based groups have been designed and studied at the B3LYP/6-311++G(3df,3pd) level of theory. The results show that $1$–$3$ are all real minima on PES, but $4$–$6$ are not. The reasons have been discussed in this work. The main conclusions of this work could be summarized as follows:

1. To design the square ptO(B)$_4$ arrangement, the strategy of "mechanical" approach is very important. Short B$_2$–C$_3$ bonds make the structures of $1$–$3$ more rigid than those of $4$–$6$, which is the key reason to achieve the square ptO(B)$_4$ arrangement successfully. The unexpected strong bonding interactions among perimeter four boron atoms and the weak interactions between O$_1$ and C$_4$ in $1$–$3$ are the other reasons to achieve the square ptO(B)$_4$ successfully.

2. The bonding interactions of the square ptO(B)$_4$ arrangement have been proposed according to the $D_{2d}$ symmetry. The shapes and symmetries of frontier orbitals, HOMO and LUMO of $1$ with $b_2$ and $e$ symmetries, respectively, have been discussed and are in accordance with the results at the B3LYP/6-311+G(3df,3pd) level of theory.

Although the $2p_\pi$ electrons of electron-rich oxygen in the square ptO(B)$_4$ arrangement are no longer the lone-pair electrons and interact with perimeter four boron atoms to form bonding MOs. The “electronic” approach, stabilized by $\pi$ acceptor substituents or by aromatic delocalization and benefited from $\sigma$ donation by electropositive groups, is also suitable for the ptO substructure.
and the Ellipticity to understand the bonding interactions of all molecules studied with ptC atoms.

The quantum theory of atoms in molecules is used here considered as strongly polar bonds. The ptO(B)4 arrangement has been studied using the B3LYP/6-311++G(3df,3pd) level of theory.

Table 2. Electron Density $\rho(r)$/(e·Bohr$^{-3}$), Laplacian of Electron Density $\nabla^2 \rho(r)$/(e·Bohr$^{-3}$), the Densities of Kinetic Energy $G(r)$/(Hartree·Bohr$^{-3}$), Potential Energy $V(r)$/(Hartree·Bohr$^{-3}$), Total Electronic Energy Density $E_\text{el}(r)$/(Hartree·Bohr$^{-3}$), and the Ellipticity $\varepsilon$ of BCPs of Some Bonds in 1–6 at the B3LYP/6-311++G(3df,3pd) Level of Theory

| molecule | bond  | $\rho(r)$ | $\nabla^2 \rho(r)$ | $E_\text{el}(r)$ | $G(r)$ | $V(r)$ | $\varepsilon$ |
|----------|-------|-----------|------------------|-----------------|-------|-------|------------|
| 1        | O1–B2 | 0.1154    | 0.5690           | −0.0756         | 0.2178 | −0.2934 | 0.4017     |
|          | B2–C3 | 0.1872    | −0.4379          | −0.1979         | 0.0884 | −0.2864 | 0.1783     |
| C3–C4    | 0.2207 | −0.4298   | −0.1620          | 0.0546          | −0.2166 | 0.0250   |
| C3–B5    | 0.1795 | −0.3164   | −0.1940          | 0.1149          | −0.3090 | 0.0967   |
| C4–C4'   | 0.2186 | −0.4273   | −0.1554          | 0.0486          | −0.2040 | 0.0307   |
| O1–O4    | 0.0305 | 0.1606    | 0.0033           | 0.0369          | −0.0336 | 4.0782   |
| 2        | O1–B2  | 0.1090    | 0.5182           | −0.0701         | 0.1996 | −0.2697 | 0.3955     |
|          | B2–C3  | 0.1901    | −0.4050          | −0.2048         | 0.1035 | −0.3083 | 0.2193     |
| C3–C4    | 0.2158 | −0.3994   | −0.1554          | 0.0555          | −0.2109 | 0.0446   |
| C3–Al5   | 0.0821 | 0.2396    | −0.0283          | 0.0881          | −0.1164 | 0.1580   |
| C4–C4'   | 0.2028 | −0.3640   | −0.1343          | 0.0432          | −0.1775 | 0.0199   |
| O1–O4    | 0.0312 | 0.1548    | 0.0029           | 0.0359          | −0.0330 | 2.6694   |
| 3        | O1–B2  | 0.1084    | 0.5233           | −0.0692         | 0.2000 | −0.2691 | 0.3607     |
|          | B2–C3  | 0.1902    | −0.3952          | −0.2059         | 0.1071 | −0.3131 | 0.2221     |
| C3–C4    | 0.2089 | −0.3623   | −0.1466          | 0.0560          | −0.2027 | 0.0711   |
| C3–Be5   | 0.1010 | 0.2312    | −0.0544          | 0.1122          | −0.1667 | 0.0924   |
| C4–C4'   | 0.1922 | −0.3221   | −0.1213          | 0.0408          | −0.1621 | 0.0153   |
| O1–O4    | 0.0307 | 0.1509    | 0.0027           | 0.0350          | −0.0323 | 2.4556   |
| 4        | O1–B2  | 0.1240    | 0.7094           | −0.0797         | 0.2571 | −0.3368 | 0.0803     |
|          | B2–C3  | 0.1869    | −0.4412          | −0.1988         | 0.0885 | −0.2873 | 0.1093     |
| C3–C4    | 0.2227 | −0.4398   | −0.1636          | 0.0537          | −0.2173 | 0.0108   |
| C3–C5    | 0.2278 | −0.4662   | −0.1731          | 0.0565          | −0.2296 | 0.0161   |
| C4–C4'   | 0.2226 | −0.4378   | −0.1616          | 0.0521          | −0.2138 | 0.0119   |
| 5        | O1–B2  | 0.1232    | 0.7050           | −0.0790         | 0.2552 | −0.3342 | 0.1292     |
|          | B2–C3  | 0.1874    | −0.4640          | −0.1980         | 0.0820 | −0.2801 | 0.0855     |
| C3–C4    | 0.2271 | −0.4646   | −0.1692          | 0.0531          | −0.2223 | 0.0243   |
| C3–N5    | 0.2445 | −0.5412   | −0.2291          | 0.0938          | −0.3230 | 0.0589   |
| C4–C4'   | 0.2253 | −0.4494   | −0.1662          | 0.0538          | −0.2201 | 0.0091   |
| 6        | O1–B2  | 0.1238    | 0.7160           | −0.0795         | 0.2571 | −0.3366 | 0.1545     |
|          | B2–C3  | 0.1886    | −0.4691          | −0.2088         | 0.0835 | −0.2843 | 0.0680     |
| C3–C4    | 0.2271 | −0.4630   | −0.1686          | 0.0528          | −0.2214 | 0.0357   |
| C3–O5    | 0.2432 | −0.5416   | −0.2910          | 0.1566          | −0.4466 | 0.0536   |
| C4–C4'   | 0.2283 | −0.4604   | −0.1713          | 0.0562          | −0.2276 | 0.0036   |
| B(OH)$_3$| –O     | 0.2183    | 0.6511           | −0.2224         | 0.3851 | −0.6075 | 0.5047     |

*Labels of atoms of 1–6 given in Figure 1.

(3) The nature of the bonding interactions in the square ptO(B)$_4$ arrangement has been studied using the quantum theory of atoms in molecules. The O1–B2 bonds with large $\pi$ bonding interactions could be considered as strongly polar bonds.

The conclusions of our work could be used to achieve the planar polycoordinate oxygen or other elements.

**COMPUTATIONAL DETAILS**

The B3LYP hybrid density functional is used in this study based on the statement that “the B3LYP hybrid density functional is still the preferred method for studying species with ptC atoms.” Geometries of all molecules studied here were optimized at the B3LYP/6-311G(d,p) level of theory first and reoptimized at the B3LYP/6-311++G(3df,3pd) level of theory. Based on the B3LYP/6-311++G(3df,3pd) geometries, properties were also calculated at the same level of theory.

The quantum theory of atoms in molecules is used here to understand the bonding interactions of all molecules studied in this work. All computations here were performed using the GAUSSIAN 09 program apart from the critical bond points and their properties of atoms in molecules, which were performed using the program Multiwf.$^{15}$

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Notes
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ABBREVIATION
ptO(X)₈ planar tetracoordinate oxygen surrounded by four X-based groups

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