Polarization dependent spontaneous-emission rate of single quantum dots in photonic crystal membranes

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(Dated: 22 December 2009)

We have measured the variation of the spontaneous emission rate with polarization for self-assembled single quantum dots in two-dimensional photonic crystal membranes. We observe a maximum anisotropy factor of 6 between the decay rates of the two bright exciton states. This large anisotropy is attributed to the substantially different projected local density of optical states for differently oriented dipoles in the photonic crystal.

PACS numbers: 78.47.da, 78.55.Cr, 78.67.Hc

Keywords: quantum dot, photonic crystal, time-resolved measurement, spontaneous emission

In the past few decades, there has been considerable interest in applying photonic crystals (PCs) for controlling the spontaneous emission (SE) of embedded emitters, which may find applications in diverse areas such as quantum information science, efficient lasers and LEDs, and for energy harvesting. Originally proposed by Yablonovitch in 1987,1 the experimental progress has been delayed due to the lack of sufficiently high quality emitters and PCs. The first experimental demonstrations of spontaneous emission control have appeared within the last five years using colloidal quantum dots or dye molecules in 3D opal PCs or quantum wells in 2D photonic crystal membranes (PCMs).2,3,4 The latter technology has proven very successful due to the excellent optical properties of self-assembled QDs, the ability to optically address single QDs, and the strongly modified optical local density of states (LDOS) in PCMs.5,6,7,8. The latter technology is considered here, is due to dark state recombination mediated by spin-flip processes.9,10 Polarization resolved spontaneous-emission measurements enable addressing each of the orthogonally polarized bright exciton states individually and thereby to probe the anisotropy of the vacuum electromagnetic field in the PCM. We quantify the polarization dependence by defining the anisotropy factor \( \eta \equiv \frac{\gamma_{X}}{\gamma_{Y}} \), where \( \gamma_{X} \) (\( \gamma_{Y} \)) represents the decay rate of the X (Y) states.

The schematic of our experimental setup is illustrated in Fig. 1. The sample is a GaAs PCM with a layer of self-assembled InAs QDs of density 250 \( \mu m^{-2} \) embedded in the center of the membrane, see Fig. 1. It is mounted in a closed-cycle cryostat at a temperature of 10 K and excited from the top by a pulsed diode laser at 780 nm (1.590 eV, which is above the bandgap of GaAs) and a repetition rate of 20 MHz. The photoluminescence (PL) is collected through a lens (NA = 0.65), sent to a monochromator, and arrives either at a CCD camera for recording emission spectra or a silicon avalanche pho-
todiode for the time-resolved measurements. In order to facilitate polarization resolved measurements, a polarizer consisting of a half-wave plate and a polarization beam-splitter is placed before the monochromator. The excitation intensity used in the measurements is about 300 mW/cm$^2$, which is below the exciton saturation level so that only photon emission from the ground state is observed. The ground state emission wavelength is centered at 950 nm (1.305 eV) with an inhomogeneous broadening of 70 meV. The resolution of the monochromator is about 120 $\mu$eV, which is larger than the energy splitting between the two bright states. However, they can still be separated by their different polarization.

During our experiments, we investigated about 30 different QDs positioned in 7 different PCMs, with the lattice parameters ranging from 260 nm to 320 nm. For the sake of exploiting a pronounced 2D PC bandgap effect, we chose QDs in PCMs with $r/a = 0.30$, where $r$ is the radius of the air holes and $a$ is the lattice constant. For comparison, we also measured decay curves of 4 QDs positioned outside the PCMs. For each QD, the PL was projected onto different polarization directions by changing the orientation of the half-wave plate before the monochromator.

Fig. 2(a) shows the PL spectrum of single QDs inside a PCM by recording either horizontal (H) or vertical (V) polarizations. The spectrum is composed of sharp emission lines originating from single QDs with linewidths limited by the resolution of the spectrometer, as indicated by the shaded area in Fig. 2(a). Fig. 2(b) displays typical decay curves for two QDs, where QD A is inside a PCM, and QD B is in the unpatterned substrate while being close in emission energy to QD A. Three decay curves for QD A are displayed corresponding to $0^\circ$ (blue, upper curve), $70^\circ$ (cyan, middle curve) or $90^\circ$ (green, lower curve) polarization. We clearly observe that the SE rate is strongly dependent on polarization illustrating that X and Y bright excitons decay significantly different in the PCM due to the anisotropic vacuum fluctuations experienced by the QD. For comparison, no such anisotropy is observed in the reference measurements on QD B. The SE rate is furthermore found to be strongly inhibited in the PCM with the inhibition factors differing for X and Y. By comparing QD A and B we derive an inhibition factor of 15.8 for the X state and 6.5 for the Y state.

Fig. 2(b) shows three decay curves for QD B (outside PCM, emission energy 1.267 eV) for $0^\circ$ (blue curve) and $90^\circ$ (green curve) polarizations that are almost on top of each other. The red lines are bi-exponential fits to the decay curves.
FIG. 3. (Color online) PL intensities and decay rates versus polarization for QD A. The triangular points (square points) are experimental results for intensities (decay rates). The solid line is the fitted result with a cosine function, and the dashed line is a guide to the eye.

expected since a strong suppression of the decay rate in the plane of the PCM results in a high emission vertically out of the membrane due to energy redistribution\(^5\). The PL intensity variation with polarization \(\theta\) is observed to follow the simple relation
\[
I = I_X + I_Y \cos(2\theta),
\]
where \(I_X\) and \(I_Y\) are the intensities of the \(X\) and \(Y\) exciton states, respectively, see Fig. 3. This can be easily understood as the result of applying polarization projection measurements on two orthogonal states.

Fig. 4 shows the anisotropy factor of decay rates for all the measured QDs measured on PCMs with various values of the lattice constant. Note that in all measurements presented in the present manuscript the QD emission was within the 2D photonic bandgap of the PCMs\(^{10}\). Large variations are observed between the individual QDs in the PCM with a maximum value of about 6. This directly demonstrates the large anisotropy of the vacuum electromagnetic field in a PC that was theoretically proposed in Ref.\(^{11}\). This anisotropy gives rise to substantial differences in the projected LDOS leading to the different decay dynamics of \(X\) and \(Y\) exciton states. For comparison, reference QDs in a bulk substrate showed no anisotropy in the decay rates for the two orthogonally polarized states.

To conclude, we have systematically measured the polarization dependent SE rate for self-assembled single QDs inside PCMs and obtained a maximum anisotropy factor of decay rate between the \(X\) and \(Y\) states of 6. Our measurement results demonstrate the large anisotropy of the vacuum electromagnetic field inside PCM\(^{10,11}\), which is a crucial condition for achieving quantum interference between two closely lying energy levels\(^{12}\) that could enable demonstration of fascinating phenomena, such as lasing without inversion\(^{16}\) or quantum beats\(^{17}\). Therefore, our experiment is not only vital in realizing complete control of the SE of single QDs with PCs, but also enables fundamental quantum optics experiments with practical systems.

ACKNOWLEDGMENTS

We thank T. Lund-Hansen and M. L. Andersen for help during the experiment, and we gratefully acknowledge financial supports from the Danish Research Council (FTP grant 274-07-0459 and FTP/FNU grant 272-09-0159).

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