Electron spin resonance measurements of a demagnetizing field on the surface of metal samples

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By comparing the signals of electron spin resonance (ESR) from two crystals of a diamond (spin-labels) the demagnetizing field of the Co, Fe, and Ni samples in the shape of strongly elongated ellipsoids of revolution (disks) has been measured. The magnetic permeabilities of the metals in the external magnetic field corresponding to the ESR of the broken chemical bonds in a natural diamond irradiated with fast reactor neutrons have been determined.

Keywords: Electron spin resonance, Diamond irradiated with neutrons, Ferromagnetic metal, Oblate ellipsoid of revolution, Magnetic permeability, Demagnetization factor.

In ESR spectroscopy, spin-labels are used for investigation of electrical and magnetic properties of materials. For example, in [1] external spin-labels were applied to measure the electrical conductivity of nonmagnetic metals and semiconductors. A plate of ruby ($\text{Al}_2\text{O}_3;\text{Cr}^{3+}$) was coated with various electric conductors and placed on a quartz holder in the center of a rectangular $H_{102}$-resonator. In this case, the ESR signal of the $\text{Cr}^{3+}$ ions was distorted and assumed the shape of the Dyson line. The ratio of the amplitudes of the low- and high-field wings of the $\text{Cr}^{3+}$ absorption line gave the impedance of the multilayer structure in the microwave field.

In [2], nickel-foil inserts placed within a microwave resonator were used to create a gradient of the polarizing magnetic field (ESR tomography method) to investigate the distribution of paramagnetic centers inside dielectric samples.

The dynamic magnetic susceptibility of antiferromagnetic CuO-based semiconductors was studied in [3] by the method of the ESR of internal spin-labels ($\text{Cu}^{2+}$ ions). Rotation of the polarization plane of microwave oscillations by the semiconductor sample in a cylindrical $H_{111}$-resonator with scanning of the polarizing magnetic field was recorded.

As far as we know, the external spin-labels have not been used for measurement of the magnetic properties of materials. The purpose of our work is to determine the demagnetizing field and the magnetic permeability of metal samples by the ESR method with the aid of two spin-labels.

We recorded the ESR spectra on a RadioPAN SE/X-2543 spectrometer in the $X$-band with automatic adjustment of the frequency of the microwave generator (klystron) to the frequency of the $H_{102}$-resonator. The polarizing magnetic field was modulated with frequencies of 100 kHz and 80 Hz; the modulation amplitude was 0.05 mT. An NMR magnetometer and a Hall probe were used to measure the magnitude of the polarizing field, and a frequency meter was used to measure the resonator frequency. The ESR signal of the ruby, positioned on a wall of the resonator, was used for monitoring the $Q$-factor of the resonator, tuning a phase of the magnetic field modulation, and calibration of the $H_1$-component of the microwave field. The ruby in the resonator was placed so that its ESR signal could not be saturated in the range of microwave powers used (up to 50 mW). The recommended volume of paramagnetic dielectric samples for such resonators is 128 mm$^3$. The measurements were carried out at room temperature.

We investigated chemically pure polycrystalline samples of Co, Fe, Ni, W, Mo, and Ta without residual magnetization. Every sample was machined to a shape close to a strongly elongated ellipsoid of revolution (hereinafter a disk) having a diameter of $d = 3$ mm and a thickness of $t = 0.1$ mm; the characteristic dimensions of the surface irregularities did not exceed 10 $\mu$m. The choice of the dimensions of the metal samples was dictated by the necessity of minimizing the absorption of the electrical component of the microwave field by them. The metal disks were oriented in the resonator so that their plane was parallel to the directions of both the external field induction $B$ and the magnetic component of the microwave field $H_1$. With this orientation, the demagnetization factor of the disk is minimal, and it can be calculated (see below).

For measurement of the demagnetizing fields, point spin-labels on the surface of ferromagnetic (Co, Fe, Ni) and nonmagnetic (W, Mo, Ta) metals were used. As the spin-labels we took two 0.4 and 0.2 mm-diameter natural diamonds of round shape irradiated by a fluence of $3 \times 10^{20}$ fast reactor neutrons/cm$^2$. The paramagnetism of the irradiated diamond is due to broken chemical bonds. The isotropic $g$-factor was determined in the usual way from the measured frequency $\nu$ of the microwave resonator and the magnetic induction $B$ of the external polarizing field [5].

To fix the samples, a cylindrical quartz rod was used in which a cut was made to provide the positioning of the metal disk and two spin-labels along the vertical axis of the $H_{102}$-resonator in the antinode of the magnetic component of the microwave field.

The larger spin-label was pasted onto the surface of the metal disk at its center and the smaller one on the quartz...
their signals. The measured values represent a single line of width 0.2 mT, which is sum of metal disks. It is evident that without the ferromagnetic field.

The sequentially recorded signals of the absorption of the microwave field by the metal disks that were alternately placed (on the quartz holder) at the center of the $H_{102}$-resonator are presented in Fig. 1. It is seen that the ferromagnetic resonance with characteristic wide lines is observed for Co, Fe, and Ni and is not manifested for W, Mo, and Ta. Note that the eigenfrequency and the $Q$-factor of the resonator change when the ferromagnetic disk is placed into the demagnetizing field leading to a decrease in the amplitude of the ESR signal. The smaller spin-label was pasted onto the quartz holder at a distance of 4.5 mm from the disk edge. The need for having a larger number of paramagnetic centers on the diamond pasted onto the metal disk is attributable to the inhomogeneity of the demagnetizing field being parallel to the sectional plane of maximum area, the demagnetization factor is 

$$H_d = N I = N \chi H_i = N (\mu_t - 1) H_i,$$  

where $N$ is the demagnetization factor, $I$ is the magnetization, $\chi$ is the magnetic susceptibility, $H_i = H - H_d$ is the magnetic field intensity in the metal, and $\mu_t(H) = 1 + \chi(H)$ is the metal relative permeability depending on the external magnetic field $H$.

For the oblate ellipsoid of revolution of diameter $d$ and thickness $t = kd$, where $k < 1$, with the external magnetic field being parallel to the sectional plane of maximum area, the demagnetization factor is [8]

$$N = \frac{k}{2(1 - k^2)^{3/2}} \left[ \frac{\pi}{2} - k(1 - k^2)^{1/2} - \arcsin k \right].$$  

FIG. 1. Signal of absorption of the microwave field of metal samples (designated by the chemical symbols of the elements) of disk shape in the magnetic field. The numerals are amplifications of the signal in recording. The polarizing field $B$ and the $H_1$-component of the microwave field are in the disk plane.

FIG. 2. ESR spectra of two diamond spin-labels. The Dia† spectrum belongs to diamonds that are 6 mm apart on the quartz holder on the plane of the disk and the broadened one is attributed to the diamond located directly on the surface of the disk, above its center.

The influence of the demagnetizing field of the disk on the diamond located on the magnetized disk is seen as displacement of its ESR line by $\delta B = \mu_0 H_d$ to the side of higher magnetic fields. The resonant values of the magnetic induction $\mu_0 H(Co) \approx 332.9$ mT, $\mu_0 H(Fe) \approx 332.7$ mT, and $\mu_0 H(Ni) \approx 332.4$ mT correspond to the center of the ESR line of the diamond which is out of the field of the ferromagnetic disk. The shift of the ESR line of the diamond on the disk to the region of higher magnetic fields with respect to that of the diamond on the quartz holder comes to $\delta B(Co) \approx 18.4$ mT, $\delta B(Fe) \approx 16.3$ mT, and $\delta B(Ni) \approx 3.4$ mT. For the W, Mo, and Ta disks, splitting of the combined line from two diamond samples is not observed. A transition from the polarizing magnetic field modulation frequency of 100 kHz to that of 80 Hz did not influence the shift $\delta B$ of the spin-label absorption line.

According to [4, 6–9], the demagnetizing field in the oblate ellipsoid of revolution made of an isotropic metal is the following:

$$H_d = N I = N \chi H_i = N (\mu_t - 1) H_i,$$  

where $N$ is the demagnetization factor, $I$ is the magnetization, $\chi$ is the magnetic susceptibility, $H_i = H - H_d$ is the magnetic field intensity in the metal, and $\mu_t(H) = 1 + \chi(H)$ is the metal relative permeability depending on the external magnetic field $H$.

For the investigated disk-shaped metal samples $k = t/d \approx 0.033$.
we have \( N \approx 0.025 \).

While eddy currents excited in the disk by the modulation field during ESR measurements can be neglected, the tangential component of the magnetic field intensity is continuous on the boundary of the ferromagnetic disk and the diamond. Then the intensity of the demagnetizing field on the disk surface \( H_d = H - H_i \), and with allowance for Eq. (1) we obtain

\[
H_d = \frac{N(\mu_r - 1)}{1 + N(\mu_r - 1)} H, \tag{3}
\]

Thus, the distance between the centers of the ESR lines of two diamond samples \( \delta B = \mu_0 H_d \), where \( H_d \) is determined from Eq. (3) with allowance for Eq. (2).

We have found experimentally that \( \delta B = \mu_0 H_d \) and, consequently, \( \mu_r(H) \) depends appreciably (especially for Fe) on the heat treatment of metals. The reason is that for the region of the ESR magnetic fields of the diamond irradiated with neutrons all the domains in the ferromagnetic disks are oriented along the external field. Moreover, the manifestation of defects of the structure of the metal samples in \( \mu_r(H) \) is not “shaded” by the influence of the reorientation of the magnetic domains, as in the region of weak magnetic fields [6–9].

Calculations of the relative permeability of the materials of the metal disks by Eq. (3) with allowance for \( \mu_0 H_d \) for Co, Fe, and Ni yield \( \mu_r(\text{Co}) \approx 3.3 \), \( \mu_r(\text{Fe}) \approx 3.1 \), and \( \mu_r(\text{Ni}) \approx 1.4 \) at an external magnetic field induction of \( B = \mu_0 H \approx 332 \text{ mT} \) and a field modulation frequency of 100 kHz.

Note that, according to (3), the value of \( N(\mu_r - 1) \) can be determined from the measured values of \( H_d = \delta B / \mu_0 \) and \( H \). Then the ratio of the relative magnetic permeability of the sample \( \mu_d \) with the demagnetization factor \( N \) to the permeability of the material of the sample \( \mu_r \) is determined from Arkad’ev’s formula [7]:

\[
\frac{\mu_d}{\mu_r} = 1 - \frac{H_d}{H} = \frac{1}{1 + N(\mu_r - 1)}.
\]

Thus, by the ESR method of spin-labels (two crystals of diamond) the demagnetizing field and magnetic permeability of the samples of the ferromagnetic metals Co, Fe, and Ni in the form of disks (strongly elongated ellipsoids of revolution) have been measured. The established possibility of experimental determination of \( \mu_r \) for the samples of canonical form [4, 6, 8], when the demagnetization factor \( N \) can be calculated, is important for both the investigation of the nature of the magnetic susceptibility of materials [10–12] and applications [13, 14].

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