1 Introduction

In the past three decades, information theoretic methods have been applied to many different systems, from molecular biology to quantum mechanical systems, and even to linguistics. Specifically speaking about quantum systems, these methods can reveal the presence of interactions, correlations of experimentally measured quantities, universal relations and much more.

The basic concept of information theory is the Shannon Information, also known as Shannon Entropy or Information Entropy. But the question arises: what exactly is information? Information, in a general sense, is whatever we get about the occurrence of a given event: for example, how surprising or unexpected results. Shannon defined [1], in 1948, an expression that measures the information (or randomness, or uncertainty, or ignorance) about a system. Obeying a set of mathematical properties defined by him (and even if subject to a certain reductionism), the information content of a system in terms of probabilities of a event to occur is:

\[ H = -K \sum_i p_i \log_b [p_i] \quad \text{or} \quad H = -K \int p(x) \log_b [p(x)] dx, \]  

respectively for the discrete and continuous cases.

From this, people started thinking about how complex a system can be and how to measure this complexity calculating it from a mathematical definition. Thus, the statistical measure of complexity introduced by Lopez-Ruiz, Mancini and Calbet [2]...
relates the complexity of a system to the information stored in it and the distance to
a situation in which all possible states of the system are equiprobable. This definition
encodes the concepts of order and disorder of a given arrangement of the system.

To illustrate these concepts we may think about two ideal systems frequently used
in physics: the ideal gas and the perfect crystal, extremes in all aspects and opposites
as well. Because they both are idealized systems, they should be thought as minimally
complex systems. However, while the latter is totally ordered, the former is totally
disordered; i.e. while for the perfect crystal one state is more probable than the
others, in the ideal gas all states are equiprobable. Summarizing this intuitive view:

- **Perfect crystal**: This system has zero complexity by definition; its strict sym-
  metry rules implies probability density centered around the prevailing state of
  perfect symmetry which result in minimal information. The system is com-
  pletely ordered.

- **Ideal gas**: Ideal gases have also zero complexity by definition; Their accessible
  states are equiprobable resulting in maximal information. The system is totally
  disordered.

The intuitive behavior/relation among complexity, information entropy and dise-
quilibrium is expected to be the one shown in Figure 1.

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![Figure 1: Intuition of what should be complexity, at least asymptotically.](image)

At this stage an suitable expression for complexity arises:
\[ C \equiv H \times D \quad \text{or} \quad C \equiv e^H \times D, \quad (2) \]

where

\[ D = \sum_i \left[ p_i - \frac{1}{N} \right]^2 \quad \text{or} \quad D = \int p^2(x) \, dx \]

is the distance to the (equilibrium) equiprobability of states.

In this work, we used the measure of complexity given by \( C = e^H D \) to study the effect of the global charge neutrality developed by Rotondo et al \cite{3} on neutron star structure following the work by de Avellar and Horvath \cite{4} who compared the complexity of sequences of stars with different equations of state using the “standard” local charge neutrality.

The standard approach to the neutron star structure assumes the relativistic hydrostatic equilibrium condition and a (realistic) equation of state. Hidden in this scheme is, by construction, the assumption of local charge neutrality leading to no global electric field inside the star. However, recent theoretical developments \cite{5} concluded that the insurgence of a critical electrical field during the gravitational collapse leads to the necessity of a full reexamination of the gravito-electrodynamical properties of neutron stars. If this is true, then one needs to consider an extension of the \( \beta \)-equilibrium condition consistently within, for example, a relativistic Thomas-Fermi equation, otherwise there could not be an equilibrium on microphysical scales.

Thus, when constructing the mathematical equations for the structure of these compact objects one needs to couple the relativistic Thomas-Fermi equation with the equilibrium condition governed by the Einstein-Maxwell equations.

In an extension of the work quoted in \cite{3}, Belvedere et al \cite{6} included the strong interaction making the coupled set of equations even more difficult to solve analytically.

Here we intend to study first the effects of the global charge neutrality in the order of the system and then to study the effects of the inclusion of the strong interaction. Our study potentially leads to the construction of a hierarchy of equations of state to be realized in nature from the informational theoretic methods.

2 Results and Conclusions

Our results show the preliminary calculations for the two density profiles from reference \cite{6} (Table \textbf{1}). It seems that the global charge neutrality and the presence of strong interactions actually lower the disequilibrium of the star sequence in a way that the star tend to the ideal gas case in our intuition plot.

If our full calculations validate this results to the entire sequence of stars, we could have a direct measure of the effects of the global charge neutrality via information
Neutrality $\rho_{\text{crust}} = 10^{10} \text{g/cm}^3$

| Density Profile | $M/M_\odot$ | $R[\text{km}]$ | $H[\text{nats}]$ | $D$ | $C$ |
|----------------|-------------|-----------------|-----------------|-----|-----|
| Global         | 2.0356      | 12.3386         | -0.586          | 30.441 | 16.947 |
| Local          | 2.2354      | 13.4787         | -0.840          | 36.611 | 15.583 |

$\rho_{\text{crust}} = 4.3 \times 10^{11} \text{g/cm}^3$

| Density Profile | $M/M_\odot$ | $R[\text{km}]$ | $H[\text{nats}]$ | $D$ | $C$ |
|----------------|-------------|-----------------|-----------------|-----|-----|
| Global         | 1.8707      | 12.5156         | -0.488          | 23.583 | 14.474 |
| Local          | 1.9794      | 13.3375         | -0.768          | 29.866 | 13.858 |

Table 1: Statistical measures of the information content, disequilibrium and complexity for two density profiles from ref. [6].

Theoretic methods on the structure of neutron stars make these interactions and conditions more probable to be realized in nature.

### 3 Perspectives and further developments

Besides the total implementation of the code to solve the structure of neutron stars with local and global charge neutrality in order to study the information content of the different equations of state, we are developing the theory to further validate the use of the density profile as a probability-like distribution to be used in the calculation of the information content of a system, avoiding the “negative” information as done so far. In particular, we defined a new density profile satisfying two important features of probability functions:

1. $p(x) \in [0:1] \ \forall \ x$,
2. $\int p(x)dx = 1$.

Using an exact density profile from a well-known exact solution of the Einstein equations we could match these conditions, yielding the results shown in Figure [?].

However, the consistency conditions alone are not enough. It is very necessary to understand better the meaning of what we want to calculate. The results for complexity shown in Figure [2] are at odds with the conclusion by de Avella and Horvath [4] and by Chatzisavvas et al. [7], who stated that “neutron stars are ordered systems that cannot grow in complexity as its mass increases”. The reason behind this difference lies in the rate at which $e^H$ increases or decreases relatively to the disequilibrium $D$ and this, in turn, is related to the signal of $H$. Thus, further studies along these lines are required to characterize how these quantities behave for self-gravitating stars.
Figure 2: Information content, disequilibrium and complexity of the sequence of stars from the exact solution of the TOV equations taking into account anisotropy in pressure and MIT Bag Model equation of state [and local charge neutrality]

References

[1] Shannon, C. E., *Bell System Technical Journal*, 27, 379-424, 623-656, 1948.

[2] López-Ruiz, R.; Mancini, H. L. and Calbet, X., *Physics Letters A*, 1995, 209, 321-326

[3] Rotondo, M., Rueda, J. A., Ruffini, R., Xue, S.-S. *Physics Letters B*, 2011, 701, 667-771

[4] de Avellar, M. G. B., Horvath, J. E., *Physics Letters A*, 2012, 376, 1085-1089

[5] Ruffini, R., Vereshchagin, G. V., Xue, S.-S., *Physics Reports*, 2010, 487, 1-4, 1-140

[6] Belvedere, R., Pugliese, D., Rueda, J. A., Ruffini, R., Xue, S.-S., *Nuclear Physics A* 2012, 883, 1-24

[7] Chatzisavvas, K. Ch.; Psonis, V. P., Panos, C. P. and Moustakidis, Ch. C., *Physics Letters A* 2009, 373, 3901-3909