Abstract. For a class of interacting particle systems in continuous space, we show that finite-volume approximations of the bulk diffusion matrix converge at an algebraic rate. The models we consider are reversible with respect to the Poisson measures with constant density, and are of non-gradient type. Our approach is inspired by recent progress in the quantitative homogenization of elliptic equations. Along the way, we develop suitable modifications of the Caccioppoli and multiscale Poincaré inequalities, which are of independent interest.

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1. Introduction

The goal of this paper is to make progress on the quantitative analysis of interacting particle systems. We consider a class of models in which each particle follows a random evolution on $\mathbb{R}^d$ which is influenced by the configuration of neighboring particles. The models we consider are reversible with respect to the Poisson measures with constant density, uniformly elliptic, and of non-gradient type. For similar models in this class, the hydrodynamic limit and the equilibrium fluctuations have been identified rigorously. In both these results, the limit object is described in terms of the bulk diffusion matrix. The main result of this paper is a proof that finite-volume approximations of this diffusion matrix converge at an algebraic rate.

Our strategy is inspired by recent developments in the quantitative analysis of elliptic equations with random coefficients, and in particular on the renormalization approach developed in [14, 13, 9, 10, 11, 7, 8]; see also [58] for a gentle introduction, and [59, 56, 38, 39, 36, 40, 37] for another approach based on concentration inequalities. This renormalization approach has shown its versatility in a number of other settings, covering now the homogenization of parabolic equations [5], finite-difference equations on percolation clusters [6, 24, 26], differential forms [25], the $\nabla \phi$ interface model [23, 12], and the Villain model [27].

Here as in the other settings mentioned above, we start from a representation of the finite-volume approximation of the bulk diffusion matrix as a family of variational problems, denoted by $\nu(U, p)$, where $U \subseteq \mathbb{R}^d$ and $p \in \mathbb{R}^d$ encodes a slope parameter. This quantity is subadditive as a function of the domain $U$. We then identify another subadditive quantity, denoted by $\nu^*(U, q)$, with $U \subseteq \mathbb{R}^d$ and $q \in \mathbb{R}^d$, such that $\nu^*(U, \cdot)$ is approximately convex dual to $\nu(U, \cdot)$. These quantities $\nu$ and $\nu^*$ provide with finite-volume lower and upper approximations of the limit diffusion matrix. Roughly speaking, the algebraic rate of convergence is obtained by showing that the defect in the convex duality between $\nu$ and $\nu^*$ can be controlled by the variation of $\nu$ and $\nu^*$ between two scales; we refer to [58, Section 3] for some intuition as to why a control of this sort is plausible.
Besides the identification of the most appropriate subadditive quantities \( \nu \) and \( \nu^* \), one of the main difficulties we encounter relates to the development of certain functional inequalities. As is to be expected, we will make use of Poincaré inequalities, which allow to control the \( L^2 \) oscillation of a function by the \( L^2 \) norm of its gradient. However, we will need to be more precise than this. Indeed, we want to be able to assert that if the gradient of a function is small in some weaker norm, then we can control the \( L^2 \) oscillation of the function more tightly. In other words, we need some analogue of the inequality \( \|u\|_{L^2} \leq C \|\nabla u\|_{H^{-1}} \). Recall that in the current paper, the functions of interest are defined over the space of all possible particle configurations. The precise statement of our “multiscale Poincaré inequality” is in Proposition 3.8.

Another crucial ingredient we need is a version of the Caccioppoli inequality. In the standard setting of elliptic equations, this inequality states that the \( L^2 \) norm of the gradient of a harmonic function can be controlled by the \( L^2 \) norm of its oscillation on a larger domain; one can think of this inequality as a “reverse Poincaré inequality” for harmonic functions. If \( u \) denotes the harmonic function, then a standard proof of this inequality consists in testing the equation for \( u \) with \( u\phi \), where \( \phi \) is a smooth cutoff function which is equal to 1 in the inner domain, and is equal to 0 outside of the larger domain.

In our context, we need to “turn off” the influence of any particle that would come too close to the boundary of the larger domain. In this case, a naive modification of the standard elliptic argument is inapplicable. This comes from the fact that, as the domains become large, there will essentially always be many particles that come dangerously close to the boundary of the larger domain; so the cutoff function \( \phi \) would essentially always have to vanish, except on an event of very small probability. We therefore need to identify a different approach. In fact, we settle for a modified form of the Caccioppoli inequality, in which we control the \( L^2 \) norm of the gradient of a solution by the \( L^2 \) norm of the solution on a larger domain, plus a fraction of the \( L^2 \) norm of its gradient on the larger domain; see Proposition 3.9 for the precise statement.

At present, we think that the results presented here should allow to derive a quantitative version of the hydrodynamic limit, as well as to derive “near-equilibrium” fluctuation results. To be more precise, for a domain of side length \( R \) and an initial density profile varying macroscopically, it should be possible to control the convergence to the hydrodynamic limit at a precision of \( R^{-\alpha} \), for some \( \alpha > 0 \). Conversely, starting from a density profile that has variations of size bounded by \( R^{-\frac{d}{2} + \alpha} \), it should be possible to identify the asymptotic fluctuations of the density field. These would represent first steps towards bridging the gap between these two results.

By analogy with the results obtained for elliptic equations and other contexts, see in particular [58, Section 3] and [11, Chapter 2 and following], we hope that the results obtained here will provide the seed for more refined, and hopefully sharp, quantitative results. This will hopefully allow to improve the exponent \( \alpha > 0 \) appearing in the previous paragraph to some explicit exponent (ideally \( \alpha = \frac{d}{2} \)), and thereby to bring us closer to a full understanding of non-equilibrium fluctuations.

We now turn to a brief overview of related works on interacting particle systems. The result in the literature that is possibly closest to ours is that of [49]. In this work, the authors consider the diffusion matrix associated with the long-time behavior of a tagged particle in the symmetric simple exclusion process, which is called the self-diffusion matrix. The main result of [49] is a proof that finite-volume approximations of the self-diffusion matrix converge to the correct limit. However, no rate of convergence could be obtained there. The qualitative result of [49] was
extended to the mean-zero simple exclusion process, and to the asymmetric simple exclusion process in dimension $d \geq 3$, in [44].

An easy consequence of the results of the present paper is that the bulk diffusion matrix is Hölder continuous as a function of the density of particles. However, for related models, it was shown in [70, 48, 15, 50, 68, 60, 61, 62] that the diffusion matrix depends smoothly on the density of particles. The situation seems comparable to that encountered when considering Bernoulli perturbations of the law of the coefficient field for elliptic equations, see [57, 31]. Possibly more difficult situations for obtaining regularity results on the homogenized parameters, with less independence built into the nature of the perturbation, include the $\nabla \phi$ model [12], and nonlinear elliptic equations [7, 8].

Two classical approaches to the identification of the hydrodynamic limit have been developed. The first, called the entropy method, was introduced in [42], and extended to certain non-gradient models in [69, 64]. The second, called the relative entropy method, was introduced in [71], and was extended to a non-gradient model in [34].

The asymptotic description of the fluctuations of interacting particle systems at equilibrium has been obtained in [18, 66, 29, 20, 22]. The extension of this result to non-gradient models was obtained in [55, 21, 33].

We are not aware of any results concerning the non-equilibrium fluctuations of a non-gradient model. For gradient models (or small perturbations thereof), we refer in particular to [63, 28, 32, 22, 45]. We also refer to the books [67, 46, 47] for much more thorough expositions on these topics, and reviews of the literature.

In relation to the purposes of the present paper, several works considered the problem of obtaining a rate of convergence to equilibrium for a system of interacting particles [54, 30, 16, 43, 51, 19, 41]. Heat kernel bounds for the tagged particle in a simple exclusion process were obtained in [35].

In all likelihood, the results presented here can be extended to other reversible models of non-gradient type, provided that the invariant measures satisfy some mixing condition (an algebraic decay of correlations would suffice, see [13]). More challenging directions include dynamics that are not uniformly elliptic, such as hard spheres. Extensions to situations in which the noise only acts on the velocity variable are likely to also be very challenging. Even further away are purely deterministic dynamics of hard spheres, as considered for instance in [17]. For any of these models, it would of course also be desirable to make progress on the quantitative analysis of the large-scale behavior of a tagged particle.

The rest of the paper is organized as follows. In Section 2, we introduce some notation and state the main result precisely, see Theorem 2.1. We then prove several functional inequalities in Section 3, including the multiscale Poincaré inequality and the modified Caccioppoli inequality. In Section 4, we define the subadditive quantities, and establish their elementary properties. Finally, in Section 5 we prove Theorem 2.1.

2. Notation and main result

In this section, we introduce some notation and state our main result.

Let $\mathcal{M}_\delta(\mathbb{R}^d)$ be the set of $\sigma$-finite measures that are sums of Dirac masses on $\mathbb{R}^d$, which we think of as the space of configurations of particles. We denote by $\mathbb{P}_\rho$ the law on $\mathcal{M}_\delta(\mathbb{R}^d)$ of the Poisson point process of density $\rho \in (0, \infty)$, with $\mathbb{E}_\rho$ the associated expectation. We denote by $\mathcal{F}_U$ the $\sigma$-algebra generated by the mappings $\mu \mapsto \mu(V)$, for all Borel sets $V \subseteq U$, completed with all the $\mathbb{P}_\rho$-null sets, and we set $\mathcal{F} := \mathcal{F}_{\mathbb{R}^d}$. 


We give ourselves a function \( a_0 : \mathcal{M}_d(\mathbb{R}^d) \to \mathbb{R}^{d \times d}_{\text{sym}} \) where \( \mathbb{R}^{d \times d}_{\text{sym}} \) is the set of \( d \)-by-\( d \) symmetric matrices. We assume that this mapping satisfies the following properties:

- **uniform ellipticity:** there exists \( \Lambda < \infty \) such that for every \( \mu \in \mathcal{M}_d(\mathbb{R}^d) \),

\[
\forall \xi \in \mathbb{R}^d, \quad |\xi|^2 \leq \xi \cdot a_0(\mu) \xi \leq \Lambda |\xi|^2 ;
\]

- **finite range of dependence:** denoting by \( B_1 \) the Euclidean ball of radius 1 centered at the origin, we assume that \( a_0 \) is \( \mathcal{F}_{B_1} \)-measurable.

We denote by \( \tau_{-x} \mu \) the translation of the measure \( \mu \) by the vector \( -x \in \mathbb{R}^d \); explicitly, for every Borel set \( U \), we have \( (\tau_{-x} \mu)(U) = \mu(x + U) \). We extend \( a_0 \) by stationarity by setting, for every \( \mu \in \mathcal{M}_d(\mathbb{R}^d) \) and \( x \in \mathbb{R}^d \),

\[
a(\mu, x) := a_0(\tau_{-x} \mu).
\]

While it would be possible to provide with a direct definition of the asymptotic bulk diffusion matrix, see for instance [46, Chapter 7], our purposes require that we identify suitable finite-volume versions of this quantity. Accordingly, for every bounded open set \( U \subseteq \mathbb{R}^d \), we define the matrix \( \bar{a}(U) \in \mathbb{R}^{d \times d}_{\text{sym}} \) to be such that, for every \( p \in \mathbb{R}^d \),

\[
(2.2) \quad \frac{1}{2} p \cdot \bar{a}(U) p = \inf_{\phi \in \mathcal{K}_0^1(U)} \mathbb{E}_\rho \left[ \frac{1}{p(U)} \int_U \frac{1}{2} (p + \nabla \phi(\mu, x)) \cdot a(\mu, x)(p + \nabla \phi(\mu, x)) \, d\mu(x) \right] .
\]

In this expression, the gradient \( \nabla \phi(\mu, x) \) is such that, for any sufficiently smooth function \( \phi \), \( x \in \text{supp} \mu \), and \( k \in \{1, \ldots, d\} \),

\[
(2.3) \quad e_k \cdot \nabla \phi(\mu, x) = \lim_{h \to 0} \frac{\phi(\mu - \delta_x + \delta_x + he_k) - \phi(\mu)}{h} ,
\]

with \( (e_1, \ldots, e_d) \) being the canonical basis of \( \mathbb{R}^d \). As will be explained in more details below, the space \( \mathcal{K}_0^1(U) \) is a completion of a space of functions that are \( \mathcal{F}_K \)-measurable for some compact set \( K \subseteq U \). The expectation \( \mathbb{E}_\rho \) is taken with respect to the variable \( \mu \), a notation we will always use to denote the canonical random variable on \( (\mathcal{M}_d(\mathbb{R}^d), \mathcal{F}, \mathbb{P}_\rho) \) (an explicit writing of \( \int_U \cdots d\mu(x) \) would actually involve a summation over every point in the intersection of \( U \) and the support of \( \mu \)).

For every \( m \in \mathbb{N} \), we let \( \square_m = Q_{3^m} \) denote the cube of side length \( 3^m \). We define the **bulk diffusion matrix** \( \bar{a} \) as

\[
\bar{a} := \lim_{m \to \infty} \bar{a}(\square_m) .
\]

Although we keep this implicit in the notation, we point out that the matrices \( \bar{a}(U) \) and \( \bar{a} \) depend on the density \( \rho \) of particles, which we keep fixed throughout the paper. The fact that this definition of \( \bar{a} \) coincides with the more classical definition, which is directly stated in infinite volume, is explained in Appendix B below. Our main result is to obtain an algebraic rate for the convergence in \( (2.4) \).

**Theorem 2.1.** The limit in \( (2.4) \) is well-defined. Moreover, there exist an exponent \( \alpha(d, \Lambda, \rho) > 0 \) and a constant \( C(d, \Lambda, \rho) < \infty \) such that for every \( m \in \mathbb{N} \),

\[
(2.5) \quad |\bar{a}(\square_m) - \bar{a}| \leq C 3^{-\alpha m} .
\]

In the remainder of this section, we clarify some of the definitions appearing earlier, and introduce some more useful notation.
2.1. Continuum configuration space. For the purposes of the present paper, we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we will not need to construct the stochastic process of interacting particles whose large-scale behavior is captured by the bulk diffusion matrix $\mathbf{a}$, so we contend ourselves with brief remarks here. Intuitively, the dynamics is a cloud of particles, which we...
The space $\mathcal{C}_c^\infty(U)$ is the subspace of $\mathcal{C}^\infty(U)$ of functions that are $\mathcal{F}_K$-measurable for some compact set $K \subseteq U$.

We now define $\mathcal{H}^1(U)$, an infinite dimensional analogue of the classical Sobolev space $H^1$. For every $f \in \mathcal{C}_c^\infty(U)$, we set

$$||f||_{\mathcal{H}^1(U)} = (\mathbb{E}_p[f^2(\mu)] + \mathbb{E}_p\left[\int_U |\nabla f(\mu, x)|^2 \, d\mu(x)\right])^{\frac{1}{2}}.$$ 

The space $\mathcal{H}^1(U)$ is the completion, with respect to this norm, of the space of functions $f \in \mathcal{C}_c^\infty(U)$ such that $||f||_{\mathcal{H}^1(U)}$ is finite (elements in this function space that coincide $\mathbb{P}_\mu$-almost surely are identified). As in classical Sobolev spaces, for every $f \in \mathcal{H}^1(U)$, we can interpret $\nabla f(\mu, x)$, with $x \in U$, in some weak sense. We stress that functions in $\mathcal{H}^1(U)$ need not be $\mathcal{F}_U$-measurable. Indeed, the function $f$ can depend on $\mu \ll U^c$ in a relatively arbitrary (measurable) way, as long as $f \in \mathcal{L}^2$.

If $V \subseteq U$ is another open set, then $\mathcal{H}^1(U) \subseteq \mathcal{H}^1(V)$.

We also define the space $\mathcal{H}^1_0(U)$ as the closure in $\mathcal{H}^1(U)$ of the space of functions $f \in \mathcal{C}_c^\infty(U)$ such that $||f||_{\mathcal{H}^1(U)}$ is finite. Notice in particular that, in stark contrast with functions in $\mathcal{H}^1(U)$, a function in $\mathcal{H}^1_0(U)$ does not depend on $\mu \ll U^c$. In the notation of Lemma 2.2, when $f \in \mathcal{H}^1_0(U)$, certain compatibility conditions between the functions $(f_n)_{n \in \mathbb{N}}$ also have to be satisfied. If $V \subseteq U$ is another open set, we have that $\mathcal{H}^1_0(V) \subseteq \mathcal{H}^1_0(U)$ (notice that the inclusion is in the opposite direction to that for $\mathcal{H}^1$ spaces). We also have the following result.

**Lemma 2.3.** For every bounded open set $U \subseteq \mathbb{R}^d$ with Lipschitz boundary and $f \in \mathcal{H}^1_0(U)$, we have

$$\mathbb{E}_\rho \left[\int_U \nabla f(\mu, x) \, d\mu(x)\right] = 0. \quad (2.6)$$

**Proof.** By density, we can assume that $f \in \mathcal{C}_c^\infty(U)$. We use the functions $(f_n)_{n \in \mathbb{N}}$ appearing in Lemma 2.2; moreover, since $f(\mu)$ does not depend on $\mu \ll U^c$, we simply write $f_n(x_1, \ldots, x_n)$ in place of $f_n(x_1, \ldots, x_n, \mu \ll U^c)$. For every $k \in \{1, \ldots, d\}$, we have

$$\mathbb{E}_\rho \left[\int_U \partial_k f(\mu, x) \, d\mu(x)\right] = \sum_{n=1}^{\infty} \mathbb{P}_\rho(\mu(\mathcal{C}) = n) \int_{U^n} e_k \cdot \nabla x_i f_n(x_1, \ldots, x_n) \, dx_1 \cdots dx_n.$$

We use Green’s formula for the integral $e_k \cdot \nabla x_i f_n(x_1, \ldots, x_n)$ with respect to $x_i$

$$\int_U e_k \cdot \nabla x_i f_n(x_1, \ldots, x_n) \, dx_i = \int_{\partial U} f_n(x_1, \ldots, x_n) e_k \cdot n(x_i) \, dx_i,$$

where $n(x_i)$ is the unit outer normal. Since $f \in \mathcal{C}_c^\infty(U)$, the quantity $f_n(x_1, \ldots, x_n)$ remains constant when $x_i$ moves along the boundary $\partial U$. Denoting this constant (which depends on $(x_j)_{j \neq i}$) by $c$, we apply once again Green’s formula to get

$$\int_U e_k \cdot \nabla x_i f_n(x_1, \ldots, x_n) \, dx_i = \int_{\partial U} c e_k \cdot n(x_i) \, dx_i = \int_U e_k \cdot \nabla x_i c \, dx_i = 0.$$

This proves the desired result. \qed

2.3. **Localization operators.** We now introduce families of operators that allow to localize a function defined on $\mathcal{M}_d(\mathbb{R}^d)$. We state some properties of these operators without proof, and refer to [41, Section 4.1] for more details.

Recall that for every $s > 0$, we write by $Q_s := \left(\frac{s}{2}, \frac{s}{2}\right)^d$. We denote the closure of the cube $Q_s$ by $\overline{Q_s}$, and define $A_s f := \mathbb{E}_\rho[f|_{\overline{Q_s}}]$. For any $f \in \mathcal{L}^2$, the process
\((A_s f)_{s \geq 0}\) is a càdlàg \(L^2\)-martingale with respect to \((\mathcal{M}_d(\mathbb{R}^d), (\mathcal{F}_{Q_s})_{s \geq 0}, \mathbb{P})\). We denote the jump at time \(s\) by
\[
\Delta_s (Af) := A_s f - A_{s-} f = A_s f - \lim_{t \searrow s} A_t f.
\]
We can have \(\Delta_s (Af) \neq 0\) only on the event where the support of the measure \(\mu\) intersects the boundary \(\partial Q_s\). The bracket process \(([Af]_s)_{s \geq 0}\) is defined by
\[
[Af]_s := \sum_{0 \leq \tau \leq s} \Delta_{\tau} (Af).
\]
We have that \(((A_s f)^2 - [Af]_s)_{s \geq 0}\) is a martingale with respect to \((\mathcal{F}_{Q_s})_{s \geq 0}\).

Notice that the operator \(A_s\) can be interpreted as an averaging of the variable \(\mu \mathbb{1}_{\overline{Q}_s}\), keeping \(\mu \mathbb{1}_{\overline{Q}_s}\) fixed. As a consequence, for every open set \(Q_s \subseteq U\), if \(f \in H^1(U)\) and \(x \in Q_s \cap \text{supp}(\mu)\), there is no ambiguity in considering the quantity \(A_s(\nabla f)(\mu, x)\). Moreover,
\[
\nabla A_s f(\mu, x) = A_s(\nabla f)(\mu, x),
\]
and \(A_s f\) belongs to \(H^1(Q_s)\), by Jensen’s inequality; see Proposition A.2 for details. However, in general, this function does not belong to \(H^1_0(\mathbb{R}^d)\), or any other \(H^1_0\) space. This comes from the fact that the function \(A_s f\) may be discontinuous as a particle enters or leaves \(Q_s\). To solve this problem, we regularize this conditional expectation in the following way. For any \(s, \varepsilon > 0\), we define
\[
A_{s, \varepsilon} f := \frac{1}{\varepsilon} \int_0^\varepsilon \mathcal{A}_{s+t} f \, dt.
\]
As above, for every open set \(U\) containing \(Q_{s+\varepsilon}\), \(f \in H^1(U)\), and \(x \in Q_{s+\varepsilon} \cap \text{supp}(\mu)\), the quantity \(A_{s, \varepsilon}(\nabla f)(\mu, x)\) is well-defined. Irrespective of the position of the point \(x \in \text{supp}(\mu)\), the gradient of \(A_{s, \varepsilon} f\) can be calculated explicitly. Indeed, writing \(\tau(x) := \inf\{r \in \mathbb{R} : x \in Q_r\}\), and \(\mathbf{n}(x)\) for the outer unit normal to \(Q_{\tau(x)}\) at the point \(x\), we have
\[
\nabla A_{s, \varepsilon} f(\mu, x) = \begin{cases} A_{s, \varepsilon}(\nabla f)(\mu, x) & \text{if } x \in Q_s; \\ \frac{1}{\varepsilon} \int_{\tau(x)-s}^\varepsilon A_{s+t}(\nabla f(\mu, x)) \, dt - \frac{\mathbf{n}(x)}{\varepsilon} \Delta_{\tau(x)}(Af) & \text{if } x \in (Q_{s+\varepsilon} \setminus Q_s); \\ 0 & \text{if } x \in Q_{s+\varepsilon}. \end{cases}
\]
Recalling that \(Q_{s+\varepsilon} \subseteq U\), one can check that \(A_{s, \varepsilon} f \in H^1_0(U)\). Similarly, one can define another regularized localization operator \(\overline{A}_{s, \varepsilon}\)
\[
\overline{A}_{s, \varepsilon} f := \frac{2}{\varepsilon^2} \int_0^\varepsilon (\varepsilon - t) A_{s+t} f \, dt,
\]
which can be obtained by applying \(A_{s, \varepsilon}\) twice: \(\overline{A}_{s, \varepsilon} = A_{s, \varepsilon} \circ A_{s, \varepsilon}\). We have the identity
\[
\mathbb{E}_\rho(\mathbb{E}_\rho(A_{s, \varepsilon} f)^2) = \mathbb{E}_\rho(A_{s, \varepsilon} f) = \mathbb{E}_\rho \left[ \frac{2}{\varepsilon^2} \int_0^\varepsilon (\varepsilon - t) (A_{s+t} f)^2 \, dt \right].
\]
The operator \(\overline{A}_{s, \varepsilon}\) satisfies properties similar to those of \(A_{s, \varepsilon}\), and we have
\[
\nabla \overline{A}_{s, \varepsilon} f(\mu, x) = \begin{cases} \overline{A}_{s, \varepsilon}(\nabla f(\mu, x)) & \text{if } x \in Q_s; \\ \frac{1}{\varepsilon} \int_{\tau(x)-s}^\varepsilon (\varepsilon - t) A_{s+t}(\nabla f(\mu, x)) \, dt - (s + \varepsilon - \tau(x)) \Delta_{\tau(x)}(Af) \mathbf{n}(x) & \text{if } x \in (Q_{s+\varepsilon} \setminus Q_s); \\ 0 & \text{if } x \in Q_{s+\varepsilon}. \end{cases}
\]
3. Functional inequalities

The goal of this section is to derive functional inequalities that will be fundamental to the proof of our main result. The first crucial estimate is a multiscale Poincaré inequality, see Proposition 3.8. This inequality is an improvement over the standard Poincaré inequality that substitutes the $L^2$ norm of the gradient of the function of interest by a weighted sum of spatial averages of this gradient. It has a structure comparable to that of $\|u\|_{L^2} \lesssim \|\nabla u\|_{H^{-1}}$, where we moreover decompose the $H^{-1}$ norm into a series a scales, in analogy with the standard definition of Besov spaces, or the equivalent definition of $H^{-1}$ norm in terms of spatial averages, see for instance [11, Appendix D]. The proof of this estimate is based on an $H^2$ estimate for solutions of $-\Delta u = f$, with $\Delta$ being the relevant Laplacian adapted to our setting; see Proposition 3.4.

The second crucial functional inequality derived here is a Caccioppoli inequality, see Proposition 3.9. In the standard elliptic setting, the Caccioppoli inequality allows to control the $L^2$ norm of the gradient of a solution by the $L^2$ norm of the function itself, on a larger domain; it can thus be thought of as a reverse Poincaré inequality for solutions. In our context, we are not able to prove such a strong estimate, but prove instead a weaker version of this inequality that allows to control the $L^2$ norm of the gradient of a solution by the $L^2$ norm of the function itself, plus a fraction of the $L^2$ norm of the gradient on a larger domain.

For every $k \leq n \in \mathbb{N}$, we define $Z_{n,k} := 3^k \mathbb{Z}^d \cap \Box_n$. Up to a set of null measure, the family $(z + \Box_n)_{z \in Z_{n,k}}$ forms a partition of $\Box_n$. For any $y \in \mathbb{R}^d$, we write $\Box_n(y)$ to denote the unique cube containing $y$ that can be written in the form $z + \Box_n$ for some $z \in 3^n \mathbb{Z}^d$. This is well-defined except for some $y$’s in a set of null measure; we can decide on an arbitrary convention for these remaining cases. We also write $Z_{n,k}(y) := 3^k \mathbb{Z}^d \cap \Box_n(y)$.

The following “multiscale spatial filtration” will be useful in the rest of the paper: for every $n, k \in \mathbb{N}$ with $k \leq n$, and $y \in \mathbb{R}^d$, we define the $\sigma$-algebra $\mathcal{G}^{y}_{n,k}$ by

$$ \mathcal{G}^{y}_{n,k} := \sigma \left( \{ \mu(z + \Box_k) \}_{z \in Z_{n,k}(y)}, \mu \ll (\mathbb{R}^d \setminus \Box_n(y)) \right). $$

We use the shorthand $\mathcal{G}_{n,k} := \mathcal{G}^{0}_{n,k}$ and $\mathcal{G}_n := \mathcal{G}_{n,n}$. One can verify that, for every $n, n', k, k' \in \mathbb{N}$ and $y, y' \in \mathbb{R}^d$,

$$ n \leq n', k \leq k' \text{ and } \Box_n(y) \subseteq \Box_{n'}(y') \implies \mathcal{G}^{y'}_{n',k'} \subseteq \mathcal{G}^{y}_{n,k}. $$

We also define the analogue of $\mathcal{G}_n$ for a general Borel set $U \subseteq \mathbb{R}^d$ as

$$ \mathcal{G}_U := \sigma \left( \mu(U), \mu \ll (\mathbb{R}^d \setminus U) \right). $$

The condition $\mathbb{E}_\rho[f|\mathcal{G}_U] = 0$ will appear many times in this paper, usually in the context of centering a function in $\mathcal{H}^1(U)$. Using the functions $(f_n)$ defined defined in Lemma 2.2, we can rewrite the condition $\mathbb{E}_\rho[f|\mathcal{G}_U] = 0$ as: for every $n \in \mathbb{N}$ and $\mathbb{P}_\rho$-almost every $\mu \in \mathcal{M}_d(\mathbb{R}^d)$,

$$ \int_{U \cap \Box_n} f_n(x_1, \cdots, x_n, \mu \ll U^c) \, dx_1 \cdots dx_n = 0. $$

3.1. Poincaré inequality. We present two types of Poincaré inequalities: one for the space $\mathcal{H}^1_0(U)$, and one for the space $\mathcal{H}^1(U)$. We first state an elementary version for product spaces and functions in the standard Sobolev $H^1$ space. The proof is classical and can be found for instance in [53, Theorem 13.36 and Proposition 13.34]. For any bounded Borel set $U \subseteq \mathbb{R}^d$, we write $\text{diam}(U)$ to denote the diameter of $U$. 


and for every \( f \in L^1(U) \), we denote the Lebesgue integral of \( f \), normalized by the Lebesgue measure of \( U \), by
\[
\int_U f := |U|^{-1} \int_U f.
\]

**Proposition 3.1** (Poincaré inequality in classical Sobolev spaces). There exists a constant \( C(d) < \infty \) such that for every bounded convex open set \( U \subseteq \mathbb{R}^d \), \( n \in \mathbb{N} \), and \( f \in H^1(U^n) \), we have
\[
\int_{U^n} \left( f - \left( \int_{U^n} f \right) \right)^2 \leq C \text{diam}(U)^2 \sum_{i=1}^n \int_{U^n} |\nabla x_i f|^2.
\]

A direct application of Proposition 3.1 gives the following proposition.

**Proposition 3.2** (Poincaré inequality in \( \mathcal{H}^1(U) \)). There exists a constant \( C(d) < \infty \) such that for every bounded convex open set and \( f \in \mathcal{H}^1(U) \), we have
\[
\mathbb{E}_\rho \left[ (f - \mathbb{E}_\rho[f|\mathcal{G}_U])^2 \right] \leq C \text{diam}(U)^2 \mathbb{E}_\rho \left[ \int_U |\nabla f|^2 \, d\mu \right].
\]

**Proof.** Without loss of generality, we may assume that \( \mathbb{E}_\rho[f|\mathcal{G}_U] = 0 \); subtracting \( \mathbb{E}_\rho[f|\mathcal{G}_U] \) from \( f \) does not change the right side of eq. (3.6). We use the functions \( (f_n) \) from Lemma 2.2, and recall that since \( \mathbb{E}_\rho[f|\mathcal{G}_U] = 0 \), we have that every function \( f_n \) is centered; see eq. (3.4). We can apply Proposition 3.1 to every \( f_n \): for a constant \( C < \infty \) independent of \( n \), we have
\[
\int_{U^n} |f_n(x_1, \ldots, x_n, \mu \cup U^c)|^2 \, dx_1 \cdots dx_n
\]
\[
\leq C \text{diam}(U)^2 \sum_{i=1}^n \int_{U^n} |\nabla x_i f_n(x_1, \ldots, x_n, \mu \cup U^c)|^2 \, dx_1 \cdots dx_n.
\]

We then sum over \( n \) and take the expectation to obtain the result. \( \square \)

Functions in the space \( \mathcal{H}^1_0(U) \) enjoy certain continuity properties as particles enter and leave the domain \( U \). For this reason, it suffices to center the function by its mean value to have a Poincaré inequality.

**Proposition 3.3** (Poincaré inequality in \( \mathcal{H}^1_0(U) \)). There exists a constant \( C(d) < \infty \) such that for every bounded open set \( U \subseteq \mathbb{R}^d \), and every \( f \in \mathcal{H}^1_0(U) \),
\[
\mathbb{E}_\rho \left[ (f - \mathbb{E}_\rho[f])^2 \right] \leq C \text{diam}(U)^2 \mathbb{E}_\rho \left[ \int_U |\nabla f|^2 \, d\mu \right].
\]

**Proof.** Without loss of generality, we assume that \( \mathbb{E}_\rho[f] = 0 \). By density, we may restrict to \( f \in \mathcal{C}_c^\infty(U) \). Applying [52, Theorem 18.7] to \( f \), we have that
\[
\mathbb{E}_\rho \left[ f^2 \right] \leq \rho \int_{\mathbb{R}^d} \mathbb{E}_\rho \left[ (f(\mu + \delta x) - f(\mu))^2 \right] \, dx.
\]

By the Fubini-Tonelli theorem, and since \( f \) is \( \mathcal{F}_U \)-measurable, this reduces to
\[
\mathbb{E}_\rho \left[ f^2 \right] \leq \rho \mathbb{E}_\rho \left[ \int_U (f(\mu + \delta x) - f(\mu))^2 \, dx \right].
\]

To establish Proposition 3.3, it thus only remains to show that
\[
\mathbb{E}_\rho \left[ \int_U (f(\mu + \delta x) - f(\mu))^2 \, dx \right] \leq \frac{C(d)}{\rho} \mathbb{E}_\rho \left[ \int_U |\nabla f|^2 \, d\mu \right].
\]
We also denote by
\( \sum_{n \in \mathbb{N}} \mathbb{P}(\mu(U) = n) \int_{U^n} \left( \int_U |f_{n+1}(x_1, \ldots, x_n, x) - f_n(x_1, \ldots, x_n)|^2 \, dx \right) dx_1 \ldots dx_n, \)
where we used the notation (similar but simpler than in Lemma 2.2)
\[
(3.10) \quad f_n(x_1, \ldots, x_n) := f \left( \sum_{k=1}^n \delta_{x_k} \right), \quad x_1, \ldots, x_n \in U.
\]

Let \( n \in \mathbb{N} \) be fixed. Since \( f \in \mathcal{C}_c^\infty(U) \), for every \( \bar{x} \in \partial U \) we have that
\[
f_n(x_1, \ldots, x_n) = f_{n+1}(x_1, \ldots, x_n, \bar{x}).
\]
That is, for every \( x_1, \ldots, x_n \in U^n \), the (smooth) function
\[
G : U \to \mathbb{R}, \quad G(\cdot) := f_{n+1}(x_1, \ldots, x_n, \cdot) - f_n(x_1, \ldots, x_n),
\]
belongs to the (standard) Sobolev space \( H^1_0(U) \). We may thus apply the standard Poincaré inequality for functions in \( H^1_0(U) \) to infer that
\[
\int_U |f_{n+1}(x_1, \ldots, x_n, x) - f_n(x_1, \ldots, x_n)|^2 \, dx \leq C(d) \text{diam}(U)^2 \int_U |\nabla_{x_1} f_{n+1}(x_1, \ldots, x_n, x)|^2 \, dx.
\]
Inserting this into (3.9), using that \( \mathbb{P}(\mu(U) = n) = e^{-n|U|/\rho(n)} \) and relabelling \( n+1 \) as \( n \), yields that
\[
\int_U \mathbb{E}_\rho \left[ (f(\cdot+\delta_x) - f(\cdot))^2 \right] \, dx \leq \frac{C(d)}{\rho} \text{diam}(U)^2 \sum_{n \in \mathbb{N}} \mathbb{P}(\mu(U) = n) n \int_{U^n} |\nabla_{x_i} f_n(x_1, \ldots, x_n)|^2 \, dx_1 \ldots dx_n.
\]
To establish (3.8) from this, it only remains to observe that, by definition (3.10) each function \( f_n \) is invariant under permutations, we have
\[
\int_{U^n} |\nabla_{x_i} f_n|^2 = \int_{U^n} |\nabla_{x_i} f_n|^2 \quad \text{for all } i = 1, \ldots, n.
\]
This concludes the proof of (3.8) and establishes Proposition 3.3. \( \square \)

3.2. \( H^2 \) estimate for the homogeneous equation. When the diffusion matrix \( a \) is a constant, the solutions to the corresponding equation have a better regularity than otherwise, and in particular, the following \( H^2 \) estimate holds. One can define the function with higher derivative iteratively: for \( x, y \in \text{supp}(\mu), x \neq y \)
\[
\partial_j \partial_k f(\mu, x, y) := \lim_{h \to 0} \frac{\partial_k f(\mu - \delta_y + \delta_{y+he_j}, x) - \partial_k f(\mu, x)}{h},
\]
and for the case \( x = y \), it makes sense as
\[
\partial_j \partial_k f(\mu, x, x) := \lim_{h \to 0} \frac{\partial_k f(\mu - \delta_x + \delta_{x+he_j}, x + he_j) - \partial_k f(\mu, x)}{h}.
\]
We also denote by \( \nabla^2 f(\mu, x, y) \) the matrix \( \{ \partial_j \partial_k f(\mu, x, y) \}_{1 \leq j, k \leq d} \), and its norm is defined as
\[
|\nabla^2 f(\mu, x, y)|^2 := \sum_{1 \leq j, k \leq d} |\partial_j \partial_k f(\mu, x, y)|^2.
\]
Proposition 3.4 ($\mathcal{H}^2$ estimate). Let $f \in L^2$, and let $u \in H^1(Q_r)$ solve “$-\Delta u = f$” in the sense that for any $v \in H^1(Q_r)$,
\begin{equation}
\mathbb{E}_\rho \left[ \int_{Q_r} \nabla u(\mu, x) \cdot \nabla v(\mu, x) \, d\mu \right] = \mathbb{E}_\rho[f v].
\end{equation}
We have the $\mathcal{H}^2(Q_r)$ estimate
\begin{equation}
\mathbb{E}_\rho \left[ \int_{(Q_r)^2} |\nabla^2 u(\mu, x, y)|^2 \, d\mu(x) d\mu(y) \right] \leq \mathbb{E}_\rho[f^2].
\end{equation}

Remark 3.5. By testing eq. (3.11) with $v = \mathbf{1}_{(\mu(Q_r) = n, \mu \in Q_r^c)}$, we see that $f$ has to satisfy $\mathbb{E}_\rho[f | G_{Q_r}] = 0$ as a condition of compatibility.

Proof of Proposition 3.4. Although this is not really part of the statement, we start by showing that for every $f \in L^2$ satisfying the compatibility condition $\mathbb{E}_\rho[f | G_{Q_r}] = 0$, there exists a solution $u$ to eq. (3.11), and we will show its link with the classical elliptic equation. At first, we notice that the problem can be studied on the space of functions
\[ W = \{ g \in H^1(Q_r) : \mathbb{E}_\rho[g | G_{Q_r}] = 0 \}. \]

Because for a general function $v \in H^1(Q_r)$, $\mathbb{E}_\rho[v | G_{Q_r}]$ can be seen as a constant in eq. (3.11): its derivative is 0 so the left-hand side of eq. (3.11) is 0. For the right-hand side, we have
\[ \mathbb{E}_\rho[f \mathbb{E}_\rho[v | G_{Q_r}]] = \mathbb{E}_\rho[\mathbb{E}_\rho[f | G_{Q_r}] \mathbb{E}_\rho[v | G_{Q_r}]] = 0. \]

Thus when applying the operation $v \mapsto v - \mathbb{E}_\rho[v | G_{Q_r}]$, we do not change eq. (3.11) and we can restrict the Laplace equation on $W$. Moreover, with the notation in Lemma 2.2, $\mathbb{E}_\rho[v | G_{Q_r}] = 0$ implies every $v_n$ is centered; see eq. (3.4).

Secondly, we test eq. (3.11) with $v \mathbf{1}_{(\mu(Q_r) = n) \cap (\mu \in Q_r^c)}$, which is conditioning the number of particles $\mu(Q_r)$, and also $(\mu \in Q_r^c(V))$ for some bounded Borel set $V$ as an environment outside $Q_r$. Then for arbitrary choices of $n, m, V$, in fact we have a classical elliptic equation using the canonical projection Lemma 2.2
\begin{equation}
\int_{(Q_r)^n} \sum_{k=1}^{n} \nabla_{x_k} u_n(x_1, \ldots, x_n, \mu \in Q_r^c) \cdot \nabla_{x_k} v_n(x_1, \ldots, x_n, \mu \in Q_r^c) \, dx_1 \cdots dx_n
\end{equation}
\[ = \int_{(Q_r)^n} f_n(x_1, \ldots, x_n, \mu \in Q_r^c) v_n(x_1, \ldots, x_n, \mu \in Q_r^c) \, dx_1 \cdots dx_n. \]

Thus the solution $u$ can be described as follows: we sample the environment outside $Q_r$ and fix the number of particle $\mu(Q_r) = n$ at first, then solve the classical elliptic equation in $H^1(\mathbb{R}^d)$ with mean zero. Finally we combine all the $u_n$ and this gives the solution of eq. (3.11). In other words, the statement of eq. (3.11) can be reinforced as
\[ \forall v \in W, \quad \mathbb{E}_\rho \left[ \int_{Q_r} \nabla u(\mu, x) \cdot \nabla v(\mu, x) \, d\mu \left| G_{Q_r} \right. \right] = \mathbb{E}_\rho \left[ f v \left| G_{Q_r} \right. \right]. \]

We now turn to study the $\mathcal{H}^2$ estimate. We apply the classical $H^2(\mathbb{R}^d)$ estimate for eq. (3.13) (see for instance [11, Lemma B.19] and its proof)
\begin{equation}
\int_{(Q_r)^n} \sum_{1 \leq i, j \leq n} |\nabla_{x_i} \nabla_{x_j} u_n|^2(x_1, \ldots, x_n, \mu \in Q_r^c) \, dx_1 \cdots dx_n
\end{equation}
\[ \leq \int_{(Q_r)^n} |f_n|^2(x_1, \ldots, x_n, \mu \in Q_r^c) \, dx_1 \cdots dx_n, \]
Taking the expectation of eq. (3.14) then gives the result. □

3.3. Multiscale Poincaré inequality. For cubes of size $3^n$, the Poincaré inequalities derived in the previous subsection (say with $k = n$ in Proposition 3.2) have a right-hand side that scales like $3^{2n}$. In this subsection, we derive a multiscale version of the Poincaré inequality, that aims to improve upon this scaling, provided that some local average of the gradient of the function is not too large. We recall that the multiscale spatial filtration $G_{n,k}^\mu$ is defined in eq. (3.1). For every $k \leq n \in \mathbb{N}$, $x, y \in \mathbb{R}^d$ such that $x \in \Box_n(y)$, open set $U$ containing $\Box_k(x)$, and $f \in \mathcal{H}^1(U)$, the following quantity is well defined

\begin{equation}
(3.15) \quad (S_{n,k}^y \nabla f)(\mu, x) := \mathbb{E}_\rho \left[ \int_{\Box_k(x)} \nabla f \, d\mu \right] G_{n,k}^y,
\end{equation}

where we use the notation, for every Borel set $V$ such that $\mu(V) \in (0, \infty)$ and function $g$ defined on $\text{supp}(\mu) \cap V$,

\begin{equation}
(3.16) \quad \int_V g \, d\mu := \frac{1}{\mu(V)} \int_V g \, d\mu,
\end{equation}

and for definiteness, we also set $\int_V g \, d\mu = 0$ if $\mu(V) = 0$. We use the shorthand notation $S_{n,k} := S_{n,k}^0$ and $S_n := S_{n,n}$. This operator has a convenient spatial martingale structure, as displayed in the following lemma.

Lemma 3.6 (Martingale structure for $S_{n,k}$). For every $n, n', k, k' \in \mathbb{N}$, $y, y' \in \mathbb{R}^d$ satisfying

\begin{equation}
n \leq n', \quad k \leq k', \quad \Box_n(y) \subseteq \Box_{n'}(y'),
\end{equation}

every $x \in \Box_{k'}(y')$, and $f \in \mathcal{H}^1(\Box_{n'}(y'))$, we have

\begin{equation}
(3.17) \quad S_{n',k'}^y \nabla f(\mu, x) = \mathbb{E}_\rho \left[ \int_{\Box_{k'}(x)} (S_{n,k}^y \nabla f) \, d\mu \right] G_{n',k'}^y.
\end{equation}

Proof. The key observation is eq. (3.2), stating that $G_{n,k}^y$ is a finer $\sigma$-algebra than $G_{n',k'}^y$, so that

\[
S_{n',k'}^y \nabla f(\mu, x)
= \mathbb{E}_\rho \left[ \frac{1}{\mu(\Box_{k'}(x))} \int_{\Box_{k'}(x)} \nabla f \, d\mu \right] G_{n',k'}^y
\]
\[
= \mathbb{E}_\rho \left[ \frac{1}{\mu(\Box_{k'}(x))} \int_{\Box_{k'}(x)} \nabla f \, d\mu \right] G_{n,k}^y G_{n',k'}^y.
\]

By the definition of $S_{n,k}^y \nabla f(\mu, z)$, we obtain

\[
S_{n',k'}^y \nabla f(\mu, x)
= \mathbb{E}_\rho \left[ \sum_{z \in \mathbb{Z}^d : \Box_{n,k} \cap \Box_{k'}(x)} \frac{\mu(z + \Box_k)}{\mu(\Box_{k'}(x))} (S_{n,k}^y \nabla f)(\mu, z) \right] G_{n',k'}^y
\]
\[
= \mathbb{E}_\rho \left[ \frac{1}{\mu(\Box_{k'}(x))} \int_{\Box_{k'}(x)} S_{n,k}^y \nabla f \, d\mu \right] G_{n',k'}^y.
\]

This is eq. (3.17). □
Moreover, for every $j,j'$, we set

$$\sum_{n,k} \prod_{x_j \in \square_k(x)} 1_{\mu(x) = \gamma_n}$$

$$\sum_{n,k} \prod_{x_j \in \square_k(x)} 1_{\mu(x) = \gamma_n}$$

Proof. Without loss of generality, we set $y = 0$. Then let $N = \sum_{z \in Z_{n,k}} N_z$ and we use the canonical projection

$$\left( S_{n,k} \nabla f \right)(\mu, x) = \frac{1}{\mu(\square_k(x))} \mathbb{E}_p \left[ \sum_{x_j \in \square_k(x)} \nabla x_j f N(\cdot, \mu \square_k) \prod_{z \in Z_{n,k}} 1_{\mu(z) = \gamma_n} G_{n,k} \right].$$

The key point is to write $\prod_{z \in Z_{n,k}} 1_{\mu(z) = \gamma_n}$ with respect to $\{z_i\}_{1 \leq i \leq N}$ such that $\mu \square_n = \sum_{i=1}^N \delta_{z_i}$. Let $\{z_i\}_{1 \leq i \leq N}$ be any fixed sequence so that every $z$ in $Z_{n,k}$ appears...
exactly $N_z$ times, as displayed in eq. (3.19). We have

$$\prod_{z \in \mathbb{Z}_{n,k}} 1_{\{\mu(z+\Box_k) = N_z\}} = \sum_{\sigma \in S_N} \prod_{i=1}^N 1_{\{x_{\sigma(i)} \in z_i + \Box_k\}},$$

where $S_N$ is the symmetric group. Moreover, under $G_{n,k}$ every permutation has equal probability, and then each $x_i$ is uniformly distributed in the associated cube $z_{\sigma(i)} + \Box_k$. Thus, we have

$$(S_{n,k} \nabla f)(\mu, x) \prod_{z \in \mathbb{Z}_{n,k}} 1_{\{\mu(z+\Box_k) = N_z\}}$$

$$= \frac{1}{\mu(\Box_k(x))} \frac{1}{|S_N|} \sum_{\sigma \in S_N} \int_{(z_{\sigma(i)} + \Box_k)_{i \in \mathbb{N}}} \nabla x_j f_N(\cdot, \mu \mathbb{L} \Box_n^c) \prod_{i=1}^N 1_{\{x_{\sigma(i)} \in z_i + \Box_k\}}.$$ 

Notice that for every $1 \leq i \leq N$, $x_i \in z_{\sigma(i)} + \Box_k$ means $x_{\sigma^{-1}(i)} \in z_i + \Box_k$, and $\sum_{x_j \in \Box_k(x)} \nabla x_j f_N(\cdot, \mu \mathbb{L} \Box_n^c)$ is permutation-invariant. So we have

$$\sum_{x_j \in \Box_k(x)} \nabla x_j f_N(x_1, \ldots, x_N, \mu \mathbb{L} \Box_n^c) \prod_{i=1}^N 1_{\{x_{\sigma(i)} \in z_i + \Box_k\}}$$

$$= \sum_{x_j \in \Box_k(x)} \nabla x_j f_N(x_1, \ldots, x_N, \mu \mathbb{L} \Box_n^c) \prod_{i=1}^N 1_{\{x_{\sigma^{-1}(i)} \in z_i + \Box_k\}}$$

$$= \sum_{x_{\sigma^{-1}(i)} \in \Box_k(x)} \nabla x_{\sigma^{-1}(i)} f_N(x_{\sigma^{-1}(1)}, \ldots, x_{\sigma^{-1}(N)}, \mu \mathbb{L} \Box_n^c) \prod_{i=1}^N 1_{\{x_{\sigma^{-1}(i)} \in z_i + \Box_k\}}$$

$$= \sum_{x_j \in \Box_k(x)} \nabla x_j f_N(x_1, \ldots, x_N, \mu \mathbb{L} \Box_n^c) \prod_{i=1}^N 1_{\{x_i \in z_i + \Box_k\}}.$$ 

Therefore, the term for each permutation has the same contribution, and we thus obtain eq. (3.18).

Then we prove eq. (3.20). To avoid possible confusion in the notation, we let $y_j, y_j'$ be the $j$-th and $j'$-th coordinates, then we exchange them and use the invariance under permutation of $f_N$,

$$e_k \cdot \nabla x_j f_N(\ldots, y_j, \ldots y_j', \ldots) = \lim_{h \to 0} \frac{f_N(\ldots, y_j + he_k, \ldots y_j', \ldots) - f_N(\ldots, y_j, \ldots y_j', \ldots)}{h}$$

(3.21)

$$= \lim_{h \to 0} \frac{f_N(\ldots, y_j', \ldots y_j + he_k, \ldots) - f_N(\ldots, y_j', \ldots y_j, \ldots)}{h} = e_k \cdot \nabla x_j f_N(\ldots, y_j', \ldots y_j, \ldots).$$

Moreover, the condition $x_j, x_j' \in \Box_k(x)$ implies that $z_j = z_j'$ and

$$1_{\{y_j \in z_j + \Box_k\}} 1_{\{y_j' \in z_j' + \Box_k\}} = 1_{\{y_j \in z_j + \Box_k\}} 1_{\{y_j' \in z_j' + \Box_k\}}.$$ 

We combine eq. (3.21) and eq. (3.22) to conclude eq. (3.20). □

We now use the operators $S_{n,k}^u$ as our locally averaged gradient to obtain the following multiscale Poincaré inequality. Notice in particular the factor of $3^k$ inside the sum on the right side of eq. (3.23), which we aim to leverage upon later by combining this with information on the smallness of $S_{n,k} \nabla u$ for $k$ close to $n$. 

\[\]
Proposition 3.8 (Multiscale Poincaré inequality). There exists a constant $C(d) < \infty$ such that for every function $u \in \mathcal{H}^1(\square_n)$ satisfying $\mathbb{E}_\rho[u | \mathcal{G}_n] = 0$, we have

\[
(3.23) \quad |u|_{\mathcal{H}^2} \leq C \left( \mathbb{E}_\rho \left[ \int_{\square_n} |\nabla u|^2 \mathrm{d}\mu \right] \right)^{\frac{1}{2}} + C \sum_{k=0}^n 3^k \left( \mathbb{E}_\rho \left[ \int_{\square_n} |S_{n,k} \nabla u|^2 \mathrm{d}\mu \right] \right)^{\frac{1}{2}}.
\]

Proof. Let $w \in \mathcal{H}^1(\square_n)$ be such that $\mathbb{E}_\rho[w | \mathcal{G}_n] = 0$ and that solves "$-\Delta w = u$", in the sense that

\[
(3.24) \quad \forall v \in \mathcal{H}^1(\square_n), \quad \mathbb{E}_\rho \left[ \int_{\square_n} \nabla w \cdot \nabla v \mathrm{d}\mu \right] = \mathbb{E}_\rho[wv],
\]

and this relation also holds conditionally on $\mathcal{G}_n$:

\[
(3.25) \quad \forall v \in \mathcal{H}^1(\square_n), \quad \mathbb{E}_\rho \left[ \int_{\square_n} \nabla w \cdot \nabla v \mathrm{d}\mu \big| \mathcal{G}_n \right] = \mathbb{E}_\rho[wv | \mathcal{G}_n].
\]

Thanks to the condition $\mathbb{E}_\rho[u | \mathcal{G}_n] = 0$, these equations are well-defined; see the proof of Proposition 3.4 for a detailed discussion. This proposition asserts that

\[
(3.26) \quad \mathbb{E}_\rho \left[ \int_{(\square_n)^2} |\nabla^2 w(\mu, x, y)|^2 \mathrm{d}\mu(x) \mathrm{d}\mu(y) \right] \leq 2C \mathbb{E}_\rho[u^2].
\]

We test eq. (3.24) with $u$ and write a telescopic sum with $(S_{n,k} \nabla w)_{0 \leq k \leq n}$ to get

\[
(3.27) \quad \mathbb{E}_\rho[u^2] = \mathbb{E}_\rho \left[ \int_{\square_n} \nabla w \cdot \nabla u \mathrm{d}\mu \right] = \text{eq. (3.27)-a + eq. (3.27)-b + eq. (3.27)-c},
\]

\[
(3.27)-a = \mathbb{E}_\rho \left[ \int_{\square_n} (\nabla w - S_{n,0} \nabla w) \cdot \nabla u \mathrm{d}\mu \right],
\]

\[
(3.27)-b = \sum_{k=0}^{n-1} \mathbb{E}_\rho \left[ \int_{\square_n} (S_{n,k} \nabla w - S_{n,k+1} \nabla w) \cdot \nabla u \mathrm{d}\mu \right],
\]

\[
(3.27)-c = \mathbb{E}_\rho \left[ \int_{\square_n} (S_{n,n} \nabla w) \cdot \nabla u \mathrm{d}\mu \right].
\]

We treat each of these three terms in turn. For eq. (3.27)-a, we use the Cauchy-Schwarz inequality to write

\[
\text{eq. (3.27)-a} \leq \left( \mathbb{E}_\rho \left[ \int_{\square_n} |\nabla w - S_{n,0} \nabla w|^2 \mathrm{d}\mu \right] \right)^{\frac{1}{2}} \left( \mathbb{E}_\rho \left[ \int_{\square_n} |\nabla u|^2 \mathrm{d}\mu \right] \right)^{\frac{1}{2}}.
\]

The first term on the right side above can be rewritten as

\[
(3.28) \quad \mathbb{E}_\rho \left[ \int_{\square_n} |\nabla w - S_{n,0} \nabla w|^2 \mathrm{d}\mu \right] = \mathbb{E}_\rho \left[ \sum_{z \in \mathbb{Z}^d_n} \mathbb{E}_\rho \left[ \int_{z+\square_0} |\nabla w - S_{n,0} \nabla w|^2 \mathrm{d}\mu \big| \mathcal{G}_{n,0} \right] \right].
\]

We use the canonical projection Lemma 2.2 for $w$ with $\mu \triangleq \sum_{i=1}^N \delta_{x_i}$, and the do the decomposition conditioned on $\mathcal{G}_{n,0}$ that

\[
w(\mu) = \sum_{N=0}^\infty \sum_{x \in \mathbb{Z}^d_n} w_N(x_1, \cdots, x_N, \mu \triangleq \sum_{i=1}^N \delta_{x_i}) \prod_{z \in \mathbb{Z}^d_n} 1_{\{|\mu(z+\square_0)| = N_z\}}.
\]

It suffices to study one term $w_N(x_1, \cdots, x_N, \mu \triangleq \sum_{i=1}^N \delta_{x_i}) \prod_{z \in \mathbb{Z}^d_n} 1_{\{|\mu(z+\square_0)| = N_z\}}$. We can apply eq. (3.18): let $\{z_i\}_{1 \leq i \leq N}$ be a fixed sequence such that eq. (3.19) holds (with...
Then we use the symmetry proved in eq. (3.20), that in fact every 
\[ \nabla (3.30) \]
\[ S \]
\[ \text{the} \]
\[ \text{Finally, we do the expectation and the sum over all} \]
\[ C \]
\[ \text{remark that the constant} \]
\[ \text{For the equation above, we can use the Poincaré inequality Proposition 3.1 because} \]
\[ \text{same contribution for all} \]
\[ 1 \]
\[ \text{ARIANNA GIUNTI, CHENLIN GU, JEAN-CHRISTOPHE MOURRAT} \]
\[ y = 0 \text{ there). For any} \ x \in \square_n \text{ we have} \]
\[ y = 0 \text{ there). For any} \ x \in \square_n \text{ we have} \]
\[ (3.29) \quad S_{n,0} \nabla w(\mu, x) \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \]
\[ = \frac{1}{\mu(\square_0(x))} \sum_{x_j \in \square_0(x)} \int_{(z + \square_0)_{1 \leq i \leq N}} \nabla x_j w_N(\cdot, \mu \square_n^c) \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}}. \]

We apply eq. (3.29) in eq. (3.28) and just study the sum over every \( z' \) in \( Z_{n,0} \)
\[ E_p \left[ \int_{z' + \square_0} |\nabla w - S_{n,0} \nabla w|^2 \, d\mu \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \bigg| \mathcal{G}_{n,0} \right] \]
\[ = \sum_{x_j \in z' + \square_0} \int_{(z + \square_0)_{1 \leq i \leq N}} \left| \nabla x_j w_N(\cdot, \mu \square_n^c) - \frac{1}{\mu(z' + \square_0)} \sum_{x_j' \in z' + \square_0} \int_{(z + \square_0)_{1 \leq i \leq N}} \nabla x_j' w_N(\cdot, \mu \square_n^c) \right|^2 \]
\[ \times \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}}. \]

Then we use the symmetry proved in eq. (3.20), that in fact every \( \nabla x_j w_N \) has the same contribution for all \( x_j \in z' + \square_0, \)
\[ E_p \left[ \int_{z' + \square_0} |\nabla w - S_{n,0} \nabla w|^2 \, d\mu \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \bigg| \mathcal{G}_{n,0} \right] \]
\[ = \sum_{x_j \in z' + \square_0} \int_{(z + \square_0)_{1 \leq i \leq N}} \left| \nabla x_j w_N(\cdot, \mu \square_n^c) - \int_{(z + \square_0)_{1 \leq i \leq N}} \nabla x_j w_N(\cdot, \mu \square_n^c) \right|^2 \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}}. \]

For the equation above, we can use the Poincaré inequality Proposition 3.1 because it is centered and every \( x_j \) lives uniformly in its associated small cube \( z_i + \square_0. \) We remark that the constant \( C \) here is independent of \( N \)
\[ E_p \left[ \int_{z' + \square_0} |\nabla w - S_{n,0} \nabla w|^2 \, d\mu \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \bigg| \mathcal{G}_{n,0} \right] \]
\[ \leq C \sum_{1 \leq i \leq N} \sum_{x_j \in z' + \square_0} \int_{(z_i + \square_0)_{1 \leq i \leq N}} \left| \nabla x_i \nabla x_j w_N(\cdot, \mu \square_n^c) \right|^2 \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}}. \]

We put this estimate back to eq. (3.28), do the sum over all \( z' \in Z_{n,0} \)
\[ \sum_{z \in Z_{n,0}} E_p \left[ \int_{z' + \square_0} |\nabla w - S_{n,0} \nabla w|^2 \, d\mu \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \bigg| \mathcal{G}_{n,0} \right] \]
\[ \leq C \sum_{1 \leq i, j \leq N} \sum_{x_j \in z' + \square_0} \int_{(z_i + \square_0)_{1 \leq i \leq N}} \left| \nabla x_i \nabla x_j w_N(\cdot, \mu \square_n^c) \right|^2 \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \]
\[ = C E_p \left[ \int_{(\square_0)^2} |\nabla^2 w(\mu, x, y)|^2 \, d\mu(x) d\mu(y) \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}} \bigg| \mathcal{G}_{n,0} \right]. \]

Finally, we do the expectation and the sum over all \( \prod_{z \in Z_{n,0}} 1_{\{\mu(z + \square_0) = N_z\}}, \) and use the \( \mathcal{H}^2 \)-estimate eq. (3.26) to obtain that
\[ E_p \left[ \int_{\square_0} |\nabla w - S_{n,0} \nabla w|^2 \, d\mu \right] \]
\[ \leq C E_p \left[ \int_{(\square_0)^2} |\nabla^2 w(\mu, x, y)|^2 \, d\mu(x) d\mu(y) \right] \]
\[ \leq C E_p [u^2]. \]
We use the definition in eq. (3.15) and Jensen’s inequality for

and this concludes that

\[
(3.31) \quad \text{eq. (3.27)} \cdot -a \leq C(\mathbb{E}_p[u^2])^{1/2} \left( \mathbb{E}_p \left[ \int_{\Box_n} |\nabla u|^2 \, d\mu \right] \right)^{1/2}.
\]

The term eq. (3.27)-b can be treated similarly. For every \( k \), we apply at first the conditional expectation with respect to \( \mathcal{G}_{n,k} \)

\[
\mathbb{E}_p \left[ \int_{\Box_n} (S_{n,k} \nabla w - S_{n,k+1} \nabla w) \cdot \nabla u \, d\mu \right]
\]

\[
= \sum_{z \in \mathcal{Z}_{n,k}} \mathbb{E}_p \left[ \mathbb{E}_p \left[ \int_{z+\Box_k} (S_{n,k} \nabla w - S_{n,k+1} \nabla w) \cdot \nabla u \, d\mu \mid \mathcal{G}_{n,k} \right] \right]
\]

\[
= \mathbb{E}_p \left[ \sum_{z \in \mathcal{Z}_{n,k}} \int_{z+\Box_k} (S_{n,k} \nabla w - S_{n,k+1} \nabla w) \cdot (S_{n,k} \nabla u) \, d\mu \right].
\]

Then we use the Cauchy-Schwarz inequality to obtain that

\[
\mathbb{E}_p \left[ \int_{\Box_n} (S_{n,k} \nabla w - S_{n,k+1} \nabla w) \cdot \nabla u \, d\mu \right]
\]

\[
\leq \left( \mathbb{E}_p \left[ \sum_{z \in \mathcal{Z}_{n,k}} \int_{z+\Box_k} |S_{n,k} \nabla w - S_{n,k+1} \nabla w|^2 \, d\mu \right] \right)^{1/2} \left( \mathbb{E}_p \left[ \sum_{z \in \mathcal{Z}_{n,k}} \int_{z+\Box_k} |S_{n,k} \nabla u|^2 \, d\mu \right] \right)^{1/2},
\]

We use the definition in eq. (3.15) and Jensen’s inequality for \(|S_{n,k} \nabla w - S_{n,k+1} \nabla w|^2\). For every \( z \in \mathcal{Z}_{n,k} \), since \((S_{n,k} \nabla w)(\mu, z)\) is \( \mathcal{G}_{n,k} \)-measurable,

\[
|S_{n,k} \nabla w - S_{n,k+1} \nabla w|^2(\mu, z) = \left( \mathbb{E}_p \left[ \int_{z+\Box_k} (\nabla w - S_{n,k+1} \nabla w) \, d\mu \mid \mathcal{G}_{n,k} \right] \right)^2
\]

\[
(3.32) \quad \leq \mathbb{E}_p \left[ \int_{z+\Box_k} |\nabla w - S_{n,k+1} \nabla w|^2 \, d\mu \mid \mathcal{G}_{n,k} \right].
\]

Then we sum over all \( z \in \mathcal{Z}_{n,k} \), and we can treat it like eq. (3.27)-a and eq. (3.30) with the Poincaré inequality in the scale \( 3^k \) and the \( \mathcal{H}^2 \)-estimate eq. (3.26), yielding

\[
\mathbb{E}_p \left[ \sum_{z \in \mathcal{Z}_{n,k}} \int_{z+\Box_k} |S_{n,k} \nabla w - S_{n,k+1} \nabla w|^2 \, d\mu \right] \leq \mathbb{E}_p \left[ \int_{\Box_n} |\nabla w - S_{n,k+1} \nabla w|^2 \, d\mu \right] \leq C 3^{2k} \mathbb{E}_p[u^2].
\]

We have thus shown that

\[
(3.33) \quad \text{eq. (3.27)} \cdot -b \leq C(\mathbb{E}_p[u^2])^{1/2} \left( \sum_{k=0}^{n-1} 3^k \left( \mathbb{E}_p \left[ \int_{\Box_n} |S_{n,k} \nabla u|^2 \, d\mu \right] \right)^{1/2} \right).
\]

For eq. (3.27)-c, we use eq. (3.15) and the Cauchy-Schwarz inequality to get that

\[
\text{eq. (3.27)} \cdot -c = \mathbb{E}_p \left[ \int_{\Box_n} (S_{n,n} \nabla w) \cdot (S_{n,n} \nabla u) \, d\mu \right]
\]

\[
\leq \left( \mathbb{E}_p \left[ \int_{\Box_n} |S_{n,n} \nabla u|^2 \, d\mu \right] \right)^{1/2} \left( \mathbb{E}_p \left[ \int_{\Box_n} |S_{n,n} \nabla u|^2 \, d\mu \right] \right)^{1/2}.
\]

To treat the term \( \mathbb{E}_p \left[ \int_{\Box_n} |S_{n,n} \nabla w|^2 \, d\mu \right] \), we define the random affine function

\[
(3.34) \quad p := \frac{(S_{n,n} \nabla w)(\mu, 0)}{|(S_{n,n} \nabla w)(\mu, 0)|}, \quad \ell_{p,\Box_n} := \int_{\Box_n} p \cdot x \, d\mu(x).
\]
We now combine eq. (3.27), (3.31), (3.33), and (3.35), to obtain eq. (3.23).

A naive approach would be to introduce a cutoff that brings the value of the function with eq. (3.25),

\[ E_\rho[u_{\ell,\square_n} | G_{n,n}] = E_\rho \left[ \int_{\square_n} \nabla w \cdot p \, d\mu \middle| G_{n,n} \right] = E_\rho \left[ \int_{\square_n} \nabla w \, d\mu \middle| G_{n,n} \right] \cdot p = \int_{\square_n} (S_{n,n} \nabla w) \cdot p \, d\mu. \]

Recalling the definition in eq. (3.34), we obtain that

\[ \int_{\square_n} |S_{n,n} \nabla w| \, d\mu = E_\rho \left[ u_{\ell,\square_n} | G_{n,n} \right] \leq (E_\rho \left[ u^2 | G_{n,n} \right])^{\frac{1}{2}} \left( E_\rho \left[ \ell^2_{p,\square_n} | G_{n,n} \right] \right)^{\frac{1}{2}} \leq C\sqrt{\mu(\square_n)} 3^n \left( E_\rho \left[ u^2 | G_{n,n} \right] \right)^{\frac{1}{2}}, \]

where in the last step, we use a direct calculation of \( E_\rho \left[ \ell^2_{p,\square_n} | G_{n,n} \right] \frac{1}{2} \), and where the constant \( C \) may depend on \( d \). Since \( S_{n,n} \nabla w \) is constant for every point in \( \square_n \), we have shown that

\[ \sqrt{\mu(\square_n)|S_{n,n} \nabla w|}(\mu,0) \leq C 3^n \left( E_\rho \left[ u^2 | G_{n,n} \right] \right)^{\frac{1}{2}}. \]

We thus obtain that

\[ E_\rho \left[ \int_{\square_n} |S_{n,n} \nabla w|^2 \, d\mu \right] = E_\rho \left[ \mu(\square_n)|S_{n,n} \nabla w|^2(\mu,0) \right] \leq C 3^{2n} E_\rho \left[ u^2 | G_{n,n} \right] = C 3^{2n} E_\rho \left[ u^2 \right], \]

and therefore

\[ (3.35) \quad \text{eq. (3.27)-c} \leq C 3^n (E_\rho[u^2])^{\frac{1}{2}} \left( E_\rho \left[ \int_{\square_n} |S_{n,n} \nabla u|^2 \, d\mu \right] \right)^{\frac{1}{2}}. \]

We now combine eq. (3.27), (3.31), (3.33), and (3.35), to obtain eq. (3.23). \( \square \)

3.4. Caccioppoli inequality. For every bounded open set \( U \subseteq \mathbb{R}^d \), we define the space of \textbf{a}-harmonic functions on \( M_\delta(\mathbb{R}^d) \) by

\[ A(U) := \left\{ u \in H^1(U) : \forall \varphi \in H^1_0(U), \quad E_\rho \left[ \int_U \nabla u \cdot a \nabla \varphi \, d\mu \right] = 0 \right\}. \]

Recalling that, for any two bounded open sets \( V \subseteq U \), we have \( H^1(U) \subseteq H^1(V) \) and \( H^1_0(V) \subseteq H^1_0(U) \), so we see that \( A(U) \subseteq A(V) \). For the classical Caccioppoli inequality, a standard proof is as follows: we multiply the harmonic function by a cutoff function, and then use this as a test function against the harmonic function itself. Adapting this argument to our space of particle configurations is not immediate. A naive approach would be to introduce a cutoff that brings the value of the function to zero whenever a particle approaches the boundary of the domain. But proceeding in this way is a very bad idea, since as we increase the size of the domain, there will essentially always be some particles near the boundary. We will instead rely on a suitable averaging procedure for particles that fall outside of a given region, using the localization operators defined in Subsection 2.3. Notice that our goal thus is not to
bring the function to zero as a particle approaches the boundary of the box. Rather, it is only to produce a function that stops depending on the position of a particle that progressively approaches the boundary of the domain, in agreement with our definition of the space \( \mathcal{H}_0^1(U) \) (and departing from the traditional definition of the Sobolev \( H_0^1 \) spaces).

**Proposition 3.9** (Modified Caccioppoli inequality). There exist \( \theta(d, \Lambda) \in (0, 1) \), \( C(d, \Lambda) < \infty \), and \( R_0(d, \Lambda) < \infty \) such that for every \( r \geq R_0 \) and \( u \in \mathcal{A}(Q_{3r}) \), we have

\[
\frac{1}{r^2 \rho(Q_3r)} \mathbb{E}_p \left[ \int_{Q_r} \nabla (A_{r,s}u) \cdot a \nabla (A_{r,s}u) \, d\mu \right] \leq \frac{C}{r^2 \rho(Q_3r)} \mathbb{E}_p [u^2] + \theta \mathbb{E}_p \left[ \frac{1}{\rho(Q_3r)} \int_{Q_{3r}} \nabla u \cdot a \nabla u \, d\mu \right].
\]

**Remark 3.10.** Inequality eq. (3.37) controls the norm of the gradient of a harmonic function in the small cube \( Q_r \), by a sum of terms involving the norm of the gradient in the larger cube \( Q_{3r} \). This does not seem to be useful at first glance. However, the key point is that the multiplicative factor \( \theta \) is smaller than one.

The proof of Proposition 3.9 will be divided into two steps. In the first step, provided by the lemma below, we prove a weaker Caccioppoli inequality, without the normalization of the volume. In the second step we use an iterative argument to improve the result and obtain Proposition 3.9.

Recall that \( A_{s,\varepsilon} \) is the regularized localization operator defined in eq. (2.9).

**Lemma 3.11** (Weak Caccioppoli inequality). Fix \( \theta'(\Lambda) := \frac{2\Lambda}{\varepsilon^{d+1}} \in (0, 1) \). For every \( r > 0, s \geq r + 2, \varepsilon > 0 \) and \( u \in \mathcal{A}(Q_{s+\varepsilon}) \), we have

\[
\frac{\theta'^2}{2 \varepsilon^2} \mathbb{E}_p [(A_{s,\varepsilon}u)^2] + \mathbb{E}_p \left[ \int_{Q_r} \nabla (A_{s,\varepsilon}u) \cdot a \nabla (A_{s,\varepsilon}u) \, d\mu \right] \leq \theta' \left( \frac{1}{2 \varepsilon^2} \mathbb{E}_p [(A_{s,\varepsilon}u)^2] + \mathbb{E}_p \left[ \int_{Q_{s+\varepsilon}} \nabla u \cdot a \nabla u \, d\mu \right] \right).
\]

**Proof.** The proof of this lemma borrows some elements from [41, Lemma 4.8]; in both settings, the main point is to construct and analyze an appropriate “cut-off” version of the function \( u \). We use the function \( \tilde{A}_{s,\varepsilon} u \in \mathcal{H}_0^1(Q_{s+\varepsilon}) \) defined in eq. (2.11) as a cut-off of the function \( u \) and test it against \( u \in \mathcal{A}(Q_{s+\varepsilon}) \) to get

\[
\mathbb{E}_p \left[ \int_{Q_{s+\varepsilon}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] = 0.
\]

Combining this with the decomposition

\[
\mathbb{E}_p \left[ \int_{Q_{s+\varepsilon}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] = \mathbb{E}_p \left[ \int_{Q_{s+2}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] + \mathbb{E}_p \left[ \int_{Q_s \cap Q_{s+2}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] + \mathbb{E}_p \left[ \int_{Q_{s+2} \setminus Q_s} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right],
\]

we have

\[
\mathbb{E}_p \left[ \int_{Q_{s+\varepsilon}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] = \mathbb{E}_p \left[ \int_{Q_{s+2}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] + \mathbb{E}_p \left[ \int_{Q_s \cap Q_{s+2}} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right] + \mathbb{E}_p \left[ \int_{Q_{s+2} \setminus Q_s} \nabla (\tilde{A}_{s,\varepsilon}u) \cdot a \nabla u \, d\mu \right],
\]
we obtain that
\begin{equation}
\text{eq. (3.40)-a} \leq |\text{eq. (3.40)-b}| + |\text{eq. (3.40)-c}|.
\end{equation}

We now study each of these three terms. For the first term eq. (3.40)-a, since \(x \in Q_{s-2}\), the coefficient \(a\) is \(\mathcal{F}_{Q_s}\)-measurable. We can thus use eq. (2.8), eq. (2.11), (2.13) and (2.12) to get
\[
\text{eq. (3.40)-a} = \frac{2}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s-2}} \int_0^\varepsilon (\varepsilon - t) A_{s+t}(\nabla u) \cdot a \nabla u \, dt \, d\mu \right] = \frac{2}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s-2}} \int_0^\varepsilon (\varepsilon - t) \mathbb{E}_{\rho} \left[ A_{s+t}(\nabla u) \cdot a A_{s+t}(\nabla u) | \mathcal{F}_{Q_{s-2}} \right] \right] \, dt \, d\mu \\
= \mathbb{E}_{\rho} \left[ \int_{Q_{s-2}} \nabla (A_{s,\varepsilon} u) \cdot a \nabla (A_{s,\varepsilon} u) \, d\mu \right].
\]

We then apply eq. (2.13) for the second term eq. (3.40)-b. We notice that in \(Q_s \setminus Q_{s-2}\), \(a\) is no longer \(\mathcal{F}_{Q_s}\)-measurable. So we use Young’s inequality and the uniform bound \(\Lambda \leq a \leq \Lambda d\)
\[
|\text{eq. (3.40)-b}| = \frac{2}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_s \setminus Q_{s-2}} \int_0^\varepsilon (\varepsilon - t) A_{s+t}(\nabla u) \cdot a \nabla u \, dt \, d\mu \right] \\
\leq \frac{\Lambda}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_s \setminus Q_{s-2}} \int_0^\varepsilon (\varepsilon - t) (|A_{s+t}(\nabla u)|^2 + |\nabla u|^2) \right] \, dt \, d\mu.
\]

For the part with conditional expectation, we use Jensen’s inequality and the uniform bound \(\Lambda \leq a \leq \Lambda d\)
\[
\frac{\Lambda}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_s \setminus Q_{s-2}} \int_0^\varepsilon \nabla u \cdot a \nabla u \, d\mu \right] \\
\leq \frac{\Lambda}{2} \mathbb{E}_{\rho} \left[ \int_{Q_s \setminus Q_{s-2}} |\nabla u|^2 \, d\mu \right] \\
\leq \frac{\Lambda}{2} \mathbb{E}_{\rho} \left[ \int_{Q_s \setminus Q_{s-2}} \nabla \cdot a \nabla u \, d\mu \right].
\]

This concludes that \(|\text{eq. (3.40)-b}| \leq \Lambda \mathbb{E}_{\rho} \left[ \int_{Q_s \setminus Q_{s-2}} \nabla u \cdot a \nabla u \, d\mu \right].

For the third term eq. (3.40)-c, we use eq. (2.13) and obtain
\[
|\text{eq. (3.40)-c}| \leq |\text{eq. (3.40)-c1}| + |\text{eq. (3.40)-c2}|
\]
\[
\text{eq. (3.40)-c1} = \frac{2}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} \int_{\tau(x)-s}^\varepsilon (\varepsilon - t) A_{s+t}(\nabla u) \cdot a \nabla u \, dt \, d\mu \right],
\]
\[
\text{eq. (3.40)-c2} = \frac{2}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} (s + \varepsilon - \tau(x)) \Delta_{\tau(x)}(Au) \nabla(x) \cdot a \nabla u \, d\mu \right].
\]

The part of eq. (3.40)-c1 can be treated as that of eq. (3.40)-b, so that
\[
\text{eq. (3.40)-c1} \leq \Lambda \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} \nabla u \cdot a \nabla u \, d\mu \right].
\]

We study the part eq. (3.40)-c2 using Young’s inequality with a parameter \(\beta > 0\) to be fixed later:
\begin{equation}
\frac{2}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} (s + \varepsilon - \tau(x)) \Delta_{\tau(x)}(Au) \nabla(x) \cdot a \nabla u \, d\mu \right] \\
\leq \frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} (s + \varepsilon - \tau(x)) |\Delta_{\tau(x)}(Au)|^2 \right] + \frac{\beta \Lambda}{\varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} (s + \varepsilon - \tau(x)) |\nabla u|^2 \, d\mu \right] \\
\leq \frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_{\rho} \left[ \int_{Q_{s+1} \setminus Q_s} (s + \varepsilon - \tau(x)) \Delta_{\tau(x)}(Au) \nabla(x) \cdot a \nabla u \, d\mu \right].
\end{equation}
The first term above will be responsible for producing the \( L^2 \) term on the right side of eq. (3.38). We start by writing

\[
\frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_\rho \left[ \int_{Q_{s+\varepsilon} \setminus Q_s} (s + \varepsilon - \tau(x)) |\Delta_{\tau(x)}(Au)|^2 \, d\mu \right] = \frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_\rho \left[ \sum_{s \leq \tau \leq s+\varepsilon} (s + \varepsilon - \tau) |\Delta_{\tau}(Au)|^2 \right],
\]

where on the right side, the sum is over all \( \tau \)'s that are jump discontinuities for \((Au)_{s>0}\). Recalling the definition of the bracket process \([Au]_s\) defined in eq. (2.7), we use Fubini’s lemma and the \( L^2 \) isometry \( \mathbb{E}_\rho \mathbb{E}[Au] = \mathbb{E}_\rho [(Au)^2] \):

\[
\frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_\rho \left[ \sum_{s \leq \tau \leq s+\varepsilon} (s + \varepsilon - \tau) |\Delta_{\tau}(Au)|^2 \right] = \frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_\rho \left[ \int_{s}^{s+\varepsilon} \sum_{s \leq \tau \leq t} |\Delta_{\tau}(Au)|^2 \, dt \right]
\]

Recalling the definition of the bracket process, we can combine this estimate with those of eq. (3.40)-a, eq. (3.40)-b, and eq. (3.40)-c1 to get

Putting this estimate back into eq. (3.42), we conclude the estimating of the term eq. (3.40)-c2, obtaining

\[
eq \frac{\Lambda}{\beta \varepsilon^2} \mathbb{E}_\rho \left[ (Au)^2 \right] - \mathbb{E}_\rho \left[ (Au)^2 \right] + \frac{\beta \Lambda}{\varepsilon} \mathbb{E}_\rho \left[ \int_{Q_{s+\varepsilon} \setminus Q_s} \nabla u \cdot a \nabla u \, d\mu \right].
\]

By choosing \( \beta = \varepsilon \), recalling eq. (3.41), and that \( r \leq s - 2 \), we can combine this estimate with those of eq. (3.40)-a, eq. (3.40)-b, and eq. (3.40)-c1 to get

\[
\frac{\Lambda}{\varepsilon^2} \mathbb{E}_\rho \left[ (Au)^2 \right] + \mathbb{E}_\rho \left[ \int_{Q_r} \nabla(Au) \cdot a \nabla (Au) \, d\mu \right] \leq \frac{\Lambda}{\varepsilon^2} \mathbb{E}_\rho \left[ (Au)^2 \right] + 2\Lambda \mathbb{E}_\rho \left[ \int_{Q_{s+\varepsilon} \setminus Q_s} \nabla u \cdot a \nabla u \, d\mu \right].
\]

We now proceed with a hole-filling argument: adding \( 2\Lambda \mathbb{E}_\rho \left[ \int_{Q_r} \nabla Au \cdot a \nabla Au \, d\mu \right] \) to both sides of the equation above, and using Jensen’s inequality, we obtain

\[
\frac{\Lambda}{\varepsilon^2} \mathbb{E}_\rho \left[ (Au)^2 \right] + (2\Lambda + 1) \mathbb{E}_\rho \left[ \int_{Q_r} \nabla(Au) \cdot a \nabla (Au) \, d\mu \right] \leq \frac{\Lambda}{\varepsilon^2} \mathbb{E}_\rho \left[ (Au)^2 \right] + 2\Lambda \mathbb{E}_\rho \left[ \int_{Q_{s+\varepsilon}} \nabla u \cdot a \nabla u \, d\mu \right].
\]

Dividing both sides by \( 2\Lambda + 1 \), and setting \( \theta' = \frac{2\Lambda}{2\Lambda+1} \), we obtain the desired inequality eq. (3.38).

We remark that eq. (3.38) does not imply directly eq. (3.37). For example, let \( r > 2 \) and we choose \( s = 2r \) and \( \varepsilon = r \) in eq. (3.38), then with a normalization of volume we
We notice that
\[
\mathbb{E}_\rho\left[ \frac{1}{\rho|Q_r|} \int_{Q_r} \nabla (A_{2r}u) \cdot a\nabla (A_{2r}u) \, d\mu \right]
\]
\[
\leq 3^d \theta' \left( \frac{1}{2r^2|Q_{3r}|} \mathbb{E}_\rho[|A_{3r}u|^2] + \mathbb{E}_\rho\left[ \frac{1}{\rho|Q_{3r}|} \int_{Q_{3r}} \nabla u \cdot a\nabla u \, d\mu \right] \right).
\]

Then another factor \(3^d\) will be added, and we typically do not have \(3^d \theta' \in (0, 1)\), since we recall that \(\theta' = \frac{2a}{2A+1}\).

**Proof of Proposition 3.9.** We apply Lemma 3.11 iteratively, with very small increments of the volume. Let \(\delta > 0\) to be fixed later, and choose \(s = (1+\delta) r, \varepsilon = \delta r\). For convenience, we assume that \(r\) is sufficiently large that
\[
(3.43) \quad s = (1+\delta) r \geq r + 2, \quad \text{that is} \quad r \geq 2\delta^{-1}.
\]

Equation (3.38) and Jensen’s inequality give us that, provided \((1+2\delta)r \leq 3r\),
\[
(3.44) \quad \mathbb{E}_\rho\left[ \frac{1}{\rho|Q_r|} \int_{Q_r} \nabla (A_{1+\delta}r, \delta u) \cdot a\nabla (A_{1+\delta}r, \delta u) \, d\mu \right]
\]
\[
\leq \tilde{\theta} \left( \frac{1}{2(2\delta r)^2|Q_{3(1+2\delta)r}|} \mathbb{E}_\rho[u^2] + \mathbb{E}_\rho\left[ \frac{1}{\rho|Q_{3(1+2\delta)r}|} \int_{Q_{3(1+2\delta)r}} \nabla u \cdot a\nabla u \, d\mu \right] \right),
\]
with \(\tilde{\theta} = (1+2\delta)^d \theta'\). We choose the constant \(\delta > 0\) sufficiently small that \(\tilde{\theta} < 1\). In order to obtain eq. (3.37), we will now apply eq. (3.44) iteratively, from the cube \(Q_r\) to the larger cube \(Q_{3r}\).

We give the details for this argument—see also Figure 2 for an illustration. We plan to use eq. (3.44) \((N+1)\) times, and and let \(\delta \in (0, 1), N \in \mathbb{N}\) satisfy
\[
(3.45) \quad \tilde{\theta} = (1+2\delta)^d \theta' < 1, \quad (1+2\delta)^{N+1} = 3.
\]

Then we set the scale and the \(a\)-harmonic functions in every scale
\[
(3.46) \quad \begin{cases} 
\frac{r_n}{n} = (1+2\delta)^n r, & 0 \leq n \leq N + 1, \\
u_{N+1} = u, & u_{N+1} = A_{1+\delta}r_n, \delta_{Q_n, u_{N+1}} \leq 0, & 0 \leq n \leq N.
\end{cases}
\]

We can prove by induction that \(u_n \in \mathcal{A}(Q_{r_n})\) under the condition eq. (3.43). Then for every \(0 \leq n \leq N\), we apply eq. (3.44) from \(u_n\) on \(Q_{r_n}\) to \(u_{n+1}\) on \(Q_{r_{n+1}}\)
\[
(3.47) \quad \mathbb{E}_\rho\left[ \frac{1}{\rho|Q_{r_n}|} \int_{Q_{r_n}} \nabla u_{n+1} \cdot a\nabla u_{n+1} \, d\mu \right]
\]
\[
\leq \tilde{\theta} \left( \frac{1}{2(2\delta r_n)^2|Q_{3(1+2\delta)r_n}|} \mathbb{E}_\rho[u_{n+1}^2] + \mathbb{E}_\rho\left[ \frac{1}{\rho|Q_{3(1+2\delta)r_n}|} \int_{Q_{3(1+2\delta)r_n}} \nabla u_{n+1} \cdot a\nabla u_{n+1} \, d\mu \right] \right).
\]

Iterating on eq. (3.47) until \(u_{N+1} = u\) on \(Q_{3r}\), we get
\[
\mathbb{E}_\rho\left[ \frac{1}{\rho|Q_{r}|} \int_{Q_{r}} \nabla u_0 \cdot a\nabla u_0 \, d\mu \right]
\]
\[
\leq \left( \frac{3^d N}{2} \sum_{n=0}^{N} (1+2\delta)^{-2n} \right) \left( \frac{1}{\rho|Q_{3r}|} \mathbb{E}[u^2] + (\tilde{\theta})^{N+1} \mathbb{E}_\rho\left[ \frac{1}{\rho|Q_{3r}|} \int_{Q_{3r}} \nabla u \cdot a\nabla u \, d\mu \right] \right).
\]
We notice that \(u_0\) can be seen as as a weighted sum of \(A_{s'}u\), for scales \(s'\) satisfying \(s' \geq (1+\delta)r \geq r + 2\), by eq. (3.43). So we apply once Jensen’s inequality for \(u_0\) and
obtain eq. (3.37) by setting

\[ C(d, \Lambda) := \frac{3^d}{2\delta^2} \sum_{n=0}^{N} (1 + 2\delta)^{-2n}, \quad \theta := (\overline{\theta})^{N+1}. \]

Although we will not use this later, we now give more explicit estimates for the choice of the parameters in the proof above, resulting from the conditions listed in eq. (3.43) and eq. (3.45). It suffices to pick an integer \( N \) larger than \( \frac{\log 3}{\log (1 + \frac{1}{2\Lambda})} \), and then in eq. (3.45) use \( \delta := \frac{1}{2}(3^{\frac{1}{N+1}} - 1) \) to fix \( \delta \), and in eq. (3.43) we require \( r \geq 2\delta^{-1} \), which gives the condition for the minimal scale \( R_0 \). A possible choice is the following

\[ N := 2 \left( \frac{d\log 3}{\log (1 + \frac{1}{2\Lambda})} \right) + 1, \quad \delta := \frac{1}{2}(3^{\frac{1}{N+1}} - 1) \approx \frac{1}{8d\Lambda}, \quad R_0 := 2\delta^{-1} \approx 16d\Lambda, \]

\( \overline{\theta} := \theta'(1 + 2\delta)^d \cdot \left( 1 + \frac{1}{2\Lambda} \right)^{-\frac{1}{2}}, \quad \theta := (\overline{\theta})^{N+1} \approx 3^{-d}, \quad C := \frac{3^d}{2\delta^2} \sum_{n=0}^{N} (1 + 2\delta)^{-2n} \approx 2^8 3^d d^3 \Lambda^3. \]

**Figure 2.** An illustration of the iterative argument for the proof of Proposition 3.9. Since Lemma 3.11 can imply Proposition 3.9 only for a comparison from scale \( r \) to \( (1 + 2\delta)r \) with \( \delta \) very small, we add many intermediate scales \( r_n = (1 + 2\delta)^n r \) between \( r \) and \( 3r \).

### 4. Subadditive Quantities

We aim to adapt the strategy in [11, Chapter 2] for our model in continuum configuration space. In this section, we define several subadditive quantities, denoted by \( \nu, \nu^*, J \), and develop their elementary properties. We then use them and a renormalization argument to obtain a quantitative rate of convergence for \( a \) in Section 5.

**4.1. Subadditive quantities \( \nu \) and \( \nu^* \).** For every bounded domain \( U \subseteq \mathbb{R}^d \) and \( p, q \in \mathbb{R}^d \), we define the affine function in \( U \) with slope \( p \) by

\[ \ell_{p,U}(\mu) := \int_{U} p \cdot x \, d\mu(x), \]
and introduce the subadditive quantities

\[
\nu(U,p) := \inf_{v \in \ell_{p,U} + \mathcal{G}_U} \mathbb{E}_p \left[ \frac{1}{\rho |U|} \int_U \frac{1}{2} \nabla v \cdot a \nabla v \, d\mu \right],
\]

(4.2)

\[
\nu^*(U,q) := \sup_{v \in \mathcal{G}_U} \mathbb{E}_p \left[ \frac{1}{\rho |U|} \int_U \left( -\frac{1}{2} \nabla u \cdot a \nabla u + q \cdot u \right) \, d\mu \right].
\]

The quantity \( \nu \) can be thought of as the average energy per unit volume of the solution which matches with the behavior of the affine function \( \ell_{p,U} \) when a particle leaves the domain \( U \). The quantity \( \nu^* \) is analogous to a Neumann problem with prescribed average flux of \( q \). As will be seen below, the quantities \( \nu \) and \( \nu^* \) are approximately dual to one another; the quality of this approximation as the domain \( U \) grows to \( \mathbb{R}^d \) will be central to the proof of Theorem 2.1. If the matrix \( a \) was constant, then by eq. (2.6) the minimizer for \( \nu(U,p) \) would be \( \ell_{p,U} \), and we would have \( \nu(U,p) = \frac{1}{2} p \cdot a p \); and similarly, were a constant, we would have \( \nu^*(U,q) = \frac{1}{2} q \cdot a^{-1} q \).

We start by recording elementary properties satisfied by \( \nu \) and \( \nu^* \). We recall that \( \mathcal{G}_U = \sigma(\mu(U), \mu_L(\mathbb{R}^d \setminus U)) \). For every \( r > 0 \), we denote by \( B_r(U) \) the \( r \)-enlargement of \( U \), that is, \( B_r(U) := \{ x \in \mathbb{R}^d : \text{dist}(x, U) < r \} \).

**Proposition 4.1** (Elementary properties of \( \nu \) and \( \nu^* \)). The following properties hold for every bounded domain \( U \subseteq \mathbb{R}^d \) with Lipschitz boundary and \( p,q,p', q' \in \mathbb{R}^d \).

1. There exists a unique solution for the optimization problem in the definition of \( \nu(U,p) \) that satisfies \( \mathbb{E}_p[\nu - \ell_{p,U}] = 0 \); we denote it by \( \nu(\cdot, U,p) \). For the optimization problem in the definition of \( \nu^*(U,q) \), there exists a maximizer \( u(\cdot, U,q) \) that is \( \mathcal{F}_{B_1(U)} \)-measurable and such that \( \mathbb{E}_p[u | \mathcal{G}_U] = 0 \). They are \( a \)-harmonic functions on \( U \), i.e. \( v(\cdot, U,p), u(\cdot, U,q) \in \mathcal{A}(U) \).

2. There exist two \( d \times d \) symmetric matrices \( a(U) \) and \( a_*(U) \) such that

\[
\nu(U,p) = \frac{1}{2} p \cdot a(U)p, \quad \nu^*(U,q) = \frac{1}{2} q \cdot a^{-1}(U)q,
\]

and these matrices satisfy \( l_d \leq a(U) \leq ld \) and \( l_d \leq a_*(U) \leq ld \). Moreover,

\[
p' \cdot a(U)p = \mathbb{E}_p \left[ \frac{1}{\rho |U|} \int_U p' \cdot a(\mu,x) \nabla v(\mu,x,U,p) \, d\mu(x) \right],
\]

(4.4)

\[
q' \cdot a^{-1}(U)q = \mathbb{E}_p \left[ \frac{1}{\rho |U|} \int_U q' \cdot \nabla u(\mu,x,U,q) \, d\mu(x) \right],
\]

(4.5)

3. Slope: \( v(\mu,U,p) \) satisfies

\[
\mathbb{E}_p \left[ \int_U \nabla v(\mu,x,U,p) \, d\mu(x) \bigg| \mathcal{G}_U \right] = \mathbb{E}_p \left[ \frac{1}{\rho |U|} \int_U \nabla v(\mu,x,U,p) \, d\mu(x) \right] = p.
\]

(4.6)

For the function \( u(\cdot, U,q) \), there exists a \( d \times d \) symmetric matrix \( l_d \leq a_*(U;G_U) \leq ld \) such that

\[
\mathbb{E}_p \left[ \int_U \nabla u(\mu,x,U,q) \, d\mu(x) \bigg| G_U \right] = a^{-1}(U;G_U)q,
\]

(4.7)

and \( a^{-1}(U) = \frac{1}{\rho |U|} \mathbb{E}_p[a^{-1}(U;G_U)\mu(U)] \), so that

\[
\mathbb{E}_p \left[ \frac{1}{\rho |U|} \int_U \nabla u(\mu,x,U,q) \, d\mu(x) \right] = a^{-1}(U)q.
\]

(4.8)
(4) Quadratic response: for every $v' \in \ell_{p,U} + \mathcal{H}^1_0(U)$, we have

$$\mathbb{E}_\rho\left[ \frac{1}{|\rho| U} \int_U \frac{1}{2} \nabla (v' - v(\mu, U, p)) \cdot a \nabla (v' - v(\mu, U, p)) \, d\mu \right] = \mathbb{E}_\rho\left[ \frac{1}{|\rho| U} \int_U \frac{1}{2} \nabla v' \cdot a \nabla v' \, d\mu \right] - \nu(U, p).$$

Similarly, for every $u' \in \mathcal{H}^1(U)$, we have

$$\mathbb{E}_\rho\left[ \frac{1}{|\rho| U} \int_U \frac{1}{2} \nabla (u' - u(\mu, U, q)) \cdot a \nabla (u' - u(\mu, U, q)) \, d\mu \right] = \nu^*(U, q) - \mathbb{E}_\rho\left[ \frac{1}{|\rho| U} \int_U \left( -\frac{1}{2} \nabla u' \cdot a \nabla u' + q \cdot \nabla u' \right) \, d\mu \right].$$

(5) The quantities $\nu$ and $\nu^*$ are subadditive: for every $n \in \mathbb{N}$,

$$\nu(\square_{n+1}, p) \leq \nu(\square_n, p), \quad \nu^*(\square_{n+1}, q) \leq \nu^*(\square_n, q).$$

Proof. We prove each of these points in turn.

(1) We study at first the maximizer for the problem $\nu^*(U, q)$. A first observation is that the maximizer can be found in $\mathcal{F}_{B_i(U)}$-measurable functions. Because for any $u \in \mathcal{H}^1(U)$, its conditional expectation $\mathbb{E}_\rho[u|\mathcal{F}_{B_i(U)}]$ reaches a larger value for the functional in $\nu^*(U, q)$. We use Jensen’s inequality that

$$\mathbb{E}_\rho\left[ \int_U \left( -\frac{1}{2} \nabla \mathbb{E}_\rho[u|\mathcal{F}_{B_i(U)}] \cdot a \nabla \mathbb{E}_\rho[u|\mathcal{F}_{B_i(U)}] + q \cdot \nabla \mathbb{E}_\rho[u|\mathcal{F}_{B_i(U)}] \right) \, d\mu \right] \geq \mathbb{E}_\rho\left[ \int_U \left( -\frac{1}{2} \nabla u \cdot a \nabla u + q \cdot \nabla u \right) \, d\mu \right].$$

By a variational calculus, we know the characterization of a maximizer with elliptic equation that for any $\phi \in \mathcal{H}^1(U)$

$$\mathbb{E}_\rho\left[ \int_U \nabla u \cdot a \nabla \phi \, d\mu \right] = \mathbb{E}_\rho\left[ \int_U q \cdot \nabla \phi \, d\mu \right].$$

Similarly to the discussion in the proof of Proposition 3.4, we know that a solution for this problem also satisfies the more precise equation

$$\mathbb{E}_\rho\left[ \int_U \nabla u \cdot a \nabla \phi \, d\mu \big| \mathcal{G}_U \right] = \mathbb{E}_\rho\left[ \int_U q \cdot \nabla \phi \, d\mu \big| \mathcal{G}_U \right],$$

and we can define its solution in the space

$$W = \{ f \in \mathcal{H}^1(U) : \mathbb{E}_\rho[ f | \mathcal{G}_U ] = 0 \}.$$  

In this space, we have

$$\mathbb{E}_\rho[f^2 | \mathcal{G}_U] \leq C \text{diam}(U)^2 \mathbb{E}_\rho\left[ \int_U |\nabla f|^2 \, d\mu \big| \mathcal{G}_U \right],$$

by the Poincaré inequality Proposition 3.2. Then the coercivity on left hand side in eq. (4.13) is ensured and we can apply the Lax-Milgram theorem. We call this maximizer $u(\mu, U, q)$. Testing eq. (4.12) with $\phi \in \mathcal{H}^1_0(U)$, eq. (2.6) implies that its right hand side is 0, so we have $u(\mu, U, q) \in \mathcal{A}(U)$.
Then we turn to $\nu(U,p)$. By a first order variation calculus, we know that a minimizer $v$ for $\nu(U,p)$ is characterized by an elliptic equation that for any $\phi \in \mathcal{H}_0^1(U)$

$$\mathbb{E}_p \left[ \int_U \nabla(v - \ell_{p,U}) \cdot a \, \nabla \phi \, d\mu \right] = \mathbb{E}_p \left[ \int_U -p \cdot a \nabla \phi \, d\mu \right].$$

(4.14)

We remark that one cannot treat this equation as eq. (4.12), because $\mathbb{E}_p[v|G_U]$ is not an element in $\mathcal{H}_0^1(U)$ and we cannot subtract it. On the other hand, we can apply the Lax-Milgram theorem on the space

$$V = \{ f \in \mathcal{H}_0^1(U) : \mathbb{E}_p[f] = 0 \},$$

to define the unique solution $v - \ell_{p,U} \in V$. We notice that the right hand side of eq. (4.14) is clearly a bounded linear functional, and the coercivity of the left hand side of eq. (4.14) is ensured by the Poincaré inequality Proposition 3.3 on $V$. We denote this minimizer by $v(\mu, U, p)$, and eq. (4.14) implies that $v(\mu, U, p) \in A(U)$.

(2) We test at first eq. (4.14) with $v(\mu, U, p') - \ell_{p',U} \in \mathcal{H}_0^1(U)$ and obtain that

$$\mathbb{E}_p \left[ \int_U \nabla(v(\mu, x, U, p)) \cdot a(\mu, x) \nabla v(\mu, x, U, p') \, d\mu(x) \right] = \mathbb{E}_p \left[ \int_U \nabla v(\mu, x, U, p) \cdot a(\mu, x) p' \, d\mu(x) \right],$$

(4.15)

and this implies $(p, p') \mapsto \mathbb{E}_p \left[ \frac{1}{|p|U} \int_U \nabla v(\mu, x, U, p) \cdot a(\mu, x) \nabla v(\mu, x, U, p') \, d\mu(x) \right]$ is a bilinear map $p \cdot \tilde{a}(U)p'$. This definition with eq. (4.15) proves eq. (4.4). We let $p = p'$ and obtain that $\nu(U, p) = \frac{1}{2} p \cdot \tilde{a}(U)p$. To obtain the bound of $\tilde{a}(U)$, we use the bound of $a$ and the definition of eq. (4.2)

$$\inf_{v \in \ell_{p,U} + \mathcal{H}_0^1(U)} \mathbb{E}_p \left[ \frac{1}{|p|U} \int_U \frac{1}{2} |\nabla v|^2 \, d\mu \right] \leq \nu(U, p) = \frac{1}{2} p \cdot \tilde{a}(U)p$$

$$\leq \inf_{v \in \ell_{p,U} + \mathcal{H}_0^1(U)} \mathbb{E}_p \left[ \frac{1}{|p|U} \int_U \frac{1}{2} |\nabla v|^2 \, d\mu \right].$$

We can check that $\ell_{p,U}$ is the minimizer for $\inf_{v \in \ell_{p,U} + \mathcal{H}_0^1(U)} \mathbb{E}_p \left[ \int_U \frac{1}{2} |\nabla v|^2 \, d\mu \right]$, then it concludes the proof of the bound $\Lambda d \leq \tilde{a}(U) \in \Lambda d$.

The same argument works for $\nu^*(U, q)$. We test eq. (4.12) with $u(\mu, U, q')$ and obtain that

$$\mathbb{E}_p \left[ \int_U \nabla u(\mu, x, U, q) \cdot a(\mu, x) \nabla u(\mu, x, U, q') \, d\mu(x) \right] = \mathbb{E}_p \left[ \int_U q \cdot \nabla u(\mu, x, U, q') \, d\mu(x) \right].$$

(4.16)

This proves that $(q, q') \mapsto \mathbb{E}_p \left[ \frac{1}{|q|U} \int_U \nabla u(\mu, x, U, q) \cdot a(\mu, x) \nabla u(\mu, x, U, q') \right]$ is also bilinear and we denote it by $q \cdot \tilde{a}_a^{-1}(U)q'$, and this also concludes eq. (4.5). Then we
put \( q' = q \) and eq. (4.16) in the definition of eq. (4.2) that

\[
\nu^*(U, q) = \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \left( -\frac{1}{2} \nabla u(\mu, x, U, q) \cdot \mathbf{a}(\mu, x) \nabla u(\mu, x, U, q) + q \cdot \nabla u(\mu, x, U, q) \right) \, d\mu(x) \right]
\]

This proves the bilinear map expression for \( \nu^*(U, q) \). Concerning the bound for the matrix \( \mathbf{a}^{-1}(U) \), we use the bound for \( \mathbf{a} \) and the equations above to obtain that

\[
\sup_{u \in \mathbb{M}^1(U)} \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \left( -\frac{\Lambda}{2} |\nabla u|^2 + q \cdot \nabla u \right) \, d\mu \right] \leq \nu^*(U, q)
\]

\[
\leq \sup_{u \in \mathbb{M}^1(U)} \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \left( \frac{1}{2} |\nabla u|^2 + q \cdot \nabla u \right) \, d\mu \right].
\]

One can check for the lower bound, \( \ell_{q, U} \) attains the maximum and for the upper bound it is \( \ell_{q, U} \) that attains the maximum. Then we put the expression \( \nu^*(U, q) = \frac{1}{2} q \cdot \mathbf{a}^{-1}(U) q \) and obtain that

\[
\frac{\Lambda^{-1}}{2} |q|^2 \leq \nu^*(U, q) = \frac{1}{2} q \cdot \mathbf{a}^{-1}(U) q \leq \frac{1}{2} |q|^2,
\]

which implies the bound for \( \mathbf{a}(U) \).

(3) The slope identity eq. (4.6) for \( v(\mu, U, p) \) is directly the result from eq. (2.6) that

\[
\mathbb{E}_\rho \left[ \int_U \nabla v(\mu, x, U, p) \, d\mu(x) \bigg| \mu(U) \right] = \mathbb{E}_\rho \left[ \int_U p \, d\mu \bigg| \mu(U) \right] = p.
\]

For the function \( u(\mu, U, q) \), the identity eq. (4.8) comes directly from eq. (4.5), but conditioned \( \mathcal{G}_U \), the averaged slope is not \( \mathbf{a}^{-1}(U) q \). In fact, we recall that \( u(\mu, U, q) \) is also the conditioned maximizer for eq. (4.13), so we can define the matrix \( \mathbf{a}^{-1}(U; \mathcal{G}_U) \), the quenched slope eq. (4.7). The estimate for this matrix is then obtained by repeating the argument in eq. (4.3) for eq. (4.13).

(4) We test eq. (4.14) with \( (v' - \ell_{p, U}) \) and put it in the left hand side of eq. (4.9)

\[
\mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla (v' - \ell_{p, U}) \cdot \mathbf{a} \nabla (v' - \ell_{p, U}) \, d\mu \right]
\]

\[
= \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla v' \cdot \mathbf{a} \nabla v' \, d\mu \right] + \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla v(\cdot, U, p) \cdot \mathbf{a} \nabla v(\cdot, U, p) \, d\mu \right]

- \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \nabla v' \cdot \mathbf{a} \nabla v(\cdot, U, p) \, d\mu \right]
\]

\[
= \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla v' \cdot \mathbf{a} \nabla v' \, d\mu \right] + \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla v(\cdot, U, p) \cdot \mathbf{a} \nabla v(\cdot, U, p) \, d\mu \right]

- \mathbb{E}_\rho \left[ \frac{1}{\rho|U|} \int_U p \cdot \mathbf{a} \nabla v(\cdot, U, p) \, d\mu \right].
\]
The term $E_{\rho} \left[ \frac{1}{\rho|U|} \int_U p \cdot a \nabla v(\cdot, U, p) \, d\mu \right]$ also appears on the right side of eq. (4.4) with $p = p'$, thus we obtain that
\[
E_{\rho} \left[ \frac{1}{\rho|U|} \int_U p \cdot a \nabla v(\cdot, U, p) \, d\mu \right] = p \cdot \tilde{a}(U) p = E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \nabla v(\cdot, U, p) \cdot a \nabla v(\cdot, U, p) \, d\mu \right],
\]
and we put it back to eq. (4.18) to conclude for the validity of eq. (4.9).

Similarly, we develop the left hand side of eq. (4.10) as eq. (4.18), and use eq. (4.12) with $\phi = u'$ to treat the inner product term of $u'$ and $u(\cdot, U, q)$.
\[
E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \nabla u' \cdot a \nabla u(\cdot, U, q) \, d\mu \right] = E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \nabla u' \cdot q \, d\mu \right].
\]
We put this term in the left hand side of eq. (4.10) and use the bilinear map expression of $\nu^*(U, q)$ to obtain that
\[
E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla (u' - u(\cdot, U, q)) \cdot a \nabla (u' - u(\cdot, U, q)) \, d\mu \right]
\]
\[
= E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla u' \cdot a \nabla u' \, d\mu \right] + E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla u(\cdot, U, q) \cdot a \nabla u(\cdot, U, q) \, d\mu \right]
\]
\[
- \mu_{\rho} \left[ \frac{1}{\rho|U|} \int_U \nabla u' \cdot a \nabla u(\cdot, U, q) \, d\mu \right]
\]
\[
= E_{\rho} \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla u' \cdot a \nabla u' \, d\mu \right] + \nu^*(U, q) - \mu_{\rho} \left[ \frac{1}{\rho|U|} \int_U \nabla u' \cdot q \, d\mu \right].
\]
This concludes the proof of eq. (4.10).

(5) For $\nu(\Box_{n+1}, p)$, we test the associated variational problem with the candidate $v' = \sum_{z \in \mathbb{Z}_{n+1, n}} v(\cdot, z + \Box_n, p)$, which is an element of $l_{p, \Box_{n+1}} + \mathcal{K}^1_0(\Box_{n+1})$, so that
\[
\nu(\Box_{n+1}, p) \leq E_{\rho} \left[ \frac{1}{\rho|\Box_{n+1}|} \int_{\Box_{n+1}} \nabla v' \cdot a \nabla v' \, d\mu \right]
\]
\[
= 3^{-d} \sum_{z \in \mathbb{Z}_{n+1, n}} E_{\rho} \left[ \frac{1}{\rho|\Box_n|} \int_{\Box_n} \nabla v(\cdot, z + \Box_n, p) \cdot a \nabla v(\cdot, z + \Box_n, p) \, d\mu \right]
\]
\[
= \nu(\Box_n, p).
\]
In the last step, we also use the stationarity of the coefficient field $a$.

For $\nu^*(\Box_{n+1}, p)$, we also use that, for every $z \in \mathbb{Z}_{n+1, n}$, we have the inclusion $\mathcal{K}^1(\Box_{n+1}) \subset \mathcal{K}^1(z + \Box_n)$, so its unit energy on every small cube $z + \Box_n$ is less than the maximum $\nu^*(z + \Box_n, p)$, thus
\[
\nu^*(\Box_{n+1}, q)
\]
\[
= 3^{-d} \sum_{z \in \mathbb{Z}_{n+1, n}} \nu^*(z + \Box_n, q)
\]
\[
= \nu^*(\Box_n, q). \quad \square
\]

4.2. Subadditive quantity $J$. We now study the quantity $J$ defined by
\[
J(U, p, q) := \nu(U, p) + \nu^*(U, q) - p \cdot q
\]
\[
= \frac{1}{2} p \cdot \tilde{a}(U) p + \frac{1}{2} q \cdot \tilde{a}^{-1}(U) q - p \cdot q.
\]
By the properties of \( \nu \) and \( \nu^* \), the quantity \( J \) is also subadditive. We briefly explain why this quantity will be convenient for our purposes. If the functions \( \nu(U, \cdot) \) and \( \nu^*(U, \cdot) \) were exactly convex dual of one another, then we would have that \( J \geq 0 \) and that for every \( p \in \mathbb{R}^d \), the infimum of \( J(U, p, \cdot) \) is zero. This would correspond to the situation in which \( \bar{\alpha}(U) \) and \( \bar{\alpha}_*(U) \) are equal, and for every \( p \in \mathbb{R}^d \), we would in fact have that \( J(U, p, \bar{\alpha}(U)p) = 0 \). Instead, we will show below that, for any symmetric matrix \( Id \leq \bar{\alpha} \leq \Lambda Id \), we have

\[
|\bar{\alpha} - \bar{\alpha}(U)| + |\bar{\alpha} - \bar{\alpha}_*(U)| \leq \sup_{p \in B_1} C(J(U, p, \bar{\alpha}p))^\frac{1}{2}.
\]

The right side of the inequality above can be thought of as a measure of the defect in the convex duality relationship between \( \nu \) and \( \nu^* \). For \( U = \square_m \) and using \( \bar{\alpha} = \bar{\alpha}_*(\square_m) \), we obtain that

\[
|\bar{\alpha}_*(\square_m) - \bar{\alpha}(\square_m)| \leq \sup_{p \in B_1} C(J(U, p, \bar{\alpha}_*(\square_m)p))^\frac{1}{2}.
\]

Since we know that \( \{\bar{\alpha}(\square_m)\}_{m \geq 0} \) is a decreasing sequence while \( \{\bar{\alpha}_*(\square_m)\}_{m \geq 0} \) is an increasing sequence from eq. (4.3) and eq. (4.11), each sequence has a limit. Therefore, once we prove a rate of convergence to zero for \( J(U, p, \bar{\alpha}_*(\square_m)p) \), we get that the two limits coincide, and also a rate for the convergence of \( \{\bar{\alpha}(\square_m)\}_{m \geq 0} \).

The rest of this section will present this strategy in details. We establish at first a variational description for the quantity \( J \) and the properties mentioned above.

**Lemma 4.2.** (1) For every \( p, q \in \mathbb{R}^d \), we have the variational representation

\[
J(U, p, q) = \sup_{w \in A(U)} E\rho\left[ \frac{1}{|\rho(U)|} \int_U \left( -\frac{1}{2} \nabla w \cdot \mathbf{a} \nabla w - p \cdot \mathbf{a} \nabla w + q \cdot \nabla w \right) \, d\mu \right].
\]

(2) We have that \( J(U, p, q) \geq 0 \) and \( \bar{\alpha}(U) \geq \bar{\alpha}_*(U) \).

(3) There exists a constant \( C(d, \Lambda) < \infty \) such that and for every symmetric matrix \( \bar{\alpha} \) satisfying \( Id \leq \bar{\alpha} \leq \Lambda Id \), we have

\[
|\bar{\alpha} - \bar{\alpha}(U)| + |\bar{\alpha} - \bar{\alpha}_*(U)| \leq C \sup_{p \in B_1} (J(U, p, \bar{\alpha}p))^\frac{1}{2}.
\]

**Proof.** (1) We start by rewriting the expression of \( J(U, p, q) \) using the definition of \( \nu^*(U, q) \) and the quadratic expression of \( \nu(U, p) \). Noting also that the maximizer of \( \nu^*(U, q) \) belongs to \( A(U) \), we can write

\[
J(U, p, q) = E\rho\left[ \frac{1}{|\rho(U)|} \int_U \frac{1}{2} \nabla v(\cdot, U, p) \cdot \mathbf{a} \nabla v(\cdot, U, p) \, d\mu \right] + \sup_{w \in A(U)} E\rho\left[ \frac{1}{|\rho(U)|} \int_U \left( -\frac{1}{2} \nabla u \cdot \mathbf{a} \nabla u + q \cdot \nabla u \right) \, d\mu \right] - p \cdot q.
\]

We claim that for any \( u \in A(U) \), with \( w := u - v(\cdot, U, p) \), we have

\[
E\rho\left[ \frac{1}{|\rho(U)|} \int_U \left( \frac{1}{2} \nabla v(\cdot, U, p) \cdot \mathbf{a} \nabla v(\cdot, U, p) - \frac{1}{2} \nabla u \cdot \mathbf{a} \nabla u + q \cdot \nabla u \right) \, d\mu \right] - p \cdot q
\]

\[
= E\rho\left[ \frac{1}{|\rho(U)|} \int_U \left( -\frac{1}{2} \nabla w \cdot \mathbf{a} \nabla w - p \cdot \mathbf{a} \nabla w + q \cdot \nabla w \right) \, d\mu \right].
\]
To prove it, we can develop the right hand side of eq. (4.23)

$$\mathbb{E}_p \left[ \frac{1}{\rho(U)} \int_U \left( -\frac{1}{2} \nabla w \cdot a \nabla w - p \cdot a \nabla w + q \cdot \nabla w \right) d\mu \right]$$

(4.24) \quad = \mathbb{E}_p \left[ \frac{1}{\rho(U)} \int_U \left( \frac{1}{2} \nabla v(\cdot, U, p) \cdot a \nabla v(\cdot, U, p) - \frac{1}{2} \nabla u \cdot a \nabla u + q \cdot \nabla u \right) d\mu \right] - p \cdot q

\quad + \mathbb{E}_p \left[ \frac{1}{\rho(U)} \int_U \nabla (v(\cdot, U, p) - \ell_p) \cdot (a \nabla u - a \nabla v(\cdot, U, p) - q) d\mu \right].

Because \((v(\cdot, U, p) - \ell_p, U) \in \mathcal{H}_0^1(U)\), we apply \(u, v(\cdot, U, p) \in \mathcal{A}(U)\) and eq. (2.6), the last line of eq. (4.24) is 0 and we prove eq. (4.23). Then we take the maximum as eq. (4.22) and obtain the definition eq. (4.20).

(2) The properties that \(J(U, p, q) \geq 0\) comes from the definition of \(\nu^*(U, q)\): we test the functional in the definition of \(\nu^*(U, q)\) with the minimizer \(v(\cdot, U, p)\) of \(\nu(U, p)\) and obtain that

\[ \nu^*(U, q) \geq \mathbb{E}_p \left[ \frac{1}{\rho(U)} \int_U \left( \frac{1}{2} \nabla v(\mu, x, U, p) \cdot a(\mu, x) \nabla v(\mu, x, U, p) + q \cdot \nabla v(\mu, x, U, p) \right) d\mu \right] = p \cdot q - \nu(U, p), \]

so that

\[ J(U, p, q) = \nu(U, p) + \nu^*(U, q) - p \cdot q \geq 0. \]

Then we test \(J(U, p, q) \geq 0\) with that \(q = \tilde{a}_*(U)p\) and obtain that

\[ 0 \leq J(U, p, \tilde{a}_*(U)p) = \frac{1}{2} p \cdot \tilde{a}(U)p + \frac{1}{2} (\tilde{a}_*(U)p) \cdot \tilde{a}_*^{-1}(U)(\tilde{a}_*(U)p) - p \cdot \tilde{a}_*(U)p, \]

and therefore \(\tilde{a}(U) \geq \tilde{a}_*(U)\).

(3) Using this property, we have

\[ J(U, p, q) = \frac{1}{2} p \cdot \tilde{a}(U)p + \frac{1}{2} q \cdot \tilde{a}_*^{-1}(U)q - p \cdot q \geq \frac{1}{2} p \cdot \tilde{a}(U)p + \frac{1}{2} q \cdot \tilde{a}_*^{-1}(U)q - p \cdot q = \frac{1}{2} (\tilde{a}(U)p - q) \tilde{a}_*^{-1}(U) \cdot (\tilde{a}(U)p - q). \]

We put \(q = \tilde{a}p\) and obtain \(|\tilde{a}(U) - \tilde{a}| \leq C \sup_{p \in B_1}(J(U, p, \tilde{a}p))^\frac{1}{2}\). The proof of the statement concerning \(|\tilde{a}_*(U) - \tilde{a}|\) is similar. \(\square\)

In view of the definition of \(J\), this functional enjoys properties similar to those described in Proposition 4.1 for \(\nu\) and \(\nu^*\).

**Proposition 4.3** (Elementary properties of \(J\)). For every bounded domain \(U \subseteq \mathbb{R}^d\) with Lipschitz boundary and \(p, q \in \mathbb{R}^d\), the quantity \(J(U, p, q)\) defined in eq. (4.19) satisfies the following properties:

(1) Characterization of optimizer: the optimization problem in eq. (4.20) admits a unique solution \(v(\cdot, U, p, q) \in \mathcal{H}_0^1(U)\) such that \(\mathbb{E}_p[v(\cdot, U, p, q)|\mathcal{G}_U] = 0\). This solution is such that for every \(w \in \mathcal{A}(U)\),

$$\mathbb{E}_p \left[ \int_U \nabla v(\cdot, U, p, q) \cdot a \nabla w d\mu \right] = \mathbb{E}_p \left[ \int_U (-p \cdot a \nabla w + q \cdot \nabla w) d\mu \right],$$

(4.25)
and \((p, q) \mapsto v(\cdot, U, p, q)\) a linear map. The function \(v(\cdot, U, p, q)\) can be expressed in terms of the optimizers in eq. (4.2) as
\[
v(\mu, U, p, q) = u(\mu, U, q) - v(\mu, U, p) - \mathbb{E}_p[u(\mu, U, q) - v(\mu, U, p)]|_{\mathcal{G}_U}.
\]
We have the quadratic expression
\[
J(U, p, q) = \mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U \frac{1}{2} \nabla v(\cdot, U, p, q) \cdot a \nabla v(\cdot, U, p, q) \, d\mu \right].
\]

(2) Slope: \(v(\cdot, U, p, q)\) satisfies
\[
\mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U \left( \frac{1}{2} \nabla (w - v(\cdot, U, p, q)) \cdot a \nabla (w - v(\cdot, U, p, q)) \right) \, d\mu \right] = J(U, p, q) - \mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U \left( \frac{1}{2} \nabla w \cdot a \nabla w - p \cdot a \nabla w + q \cdot \nabla w \right) \, d\mu \right].
\]

(3) Quadratic response: for every \(w \in \mathcal{A}(U)\), we have
\[
\mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U \nabla v(\cdot, U, p_1 + p_2, q_1 + q_2) \cdot a \nabla w \, d\mu \right] = \mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U - (p_1 + p_2) \cdot a \nabla w + (q_1 + q_2) \cdot \nabla w \, d\mu \right] = \mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U \nabla v(\cdot, U, p_1, q_1) + v(\cdot, U, p_2, q_2) \cdot a \nabla w \, d\mu \right].
\]

Then \((v(\cdot, U, p_1, q_1) + v(\cdot, U, p_2, q_2))\) is also a solution for the problem eq. (4.25) with parameter \((p_1 + p_2, q_1 + q_2)\). Notice that we have
\[
\mathbb{E}_p[(v(\cdot, U, p_1, q_1) + v(\cdot, U, p_2, q_2)) |_{\mathcal{G}_U}] = 0,
\]
it implies \(v(\mu, U, p_1 + p_2, q_1 + q_2) = v(\mu, U, p_1, q_1) + v(\mu, U, p_2, q_2)\) and the linearity of the map.

The exact expression of \(v(\mu, U, P, q)\) comes from the equivalent definition eq. (4.20) of \(J(U, p, q)\) and its proof. We put \(v(\mu, U, p, q)\) in the first order variation eq. (4.25)
\[
\mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U (-p \cdot a \nabla v(\cdot, U, p, q) + q \cdot \nabla v(\cdot, U, p, q)) \, d\mu \right] = \mathbb{E}_p \left[ \frac{1}{\rho|U|} \int_U \nabla v(\cdot, U, p, q) \cdot a \nabla v(\cdot, U, p, q) \, d\mu \right].
\]
Then we put this equation into eq. (4.20) to get eq. (4.27).
Then we add back the term $\nu$.

Lemma 4.4 (Comparison between two scales)

(4.31) 1

\[ J_{eq. (4.29)} \text{ for } p, q \in \mathbb{R}^d \]

We conclude this section with the following lemma.

**Lemma 4.4** (Comparison between two scales). For every $n, k \in \mathbb{N}$ with $k \leq n$, and $p, q \in \mathbb{R}^d$, writing $v(U)$ as shorthand for $v(\cdot, U, p, q)$, we have

\[
\frac{1}{|Z_{n, k}|} \sum_{z \in Z_{n, k}} \mathbb{E}_\rho \left[ \frac{1}{\rho|\square|} \int_{z+\square} \frac{1}{2} |\nabla v(\square_n) - \nabla v(z + \square_k)|^2 \, d\mu \right] \leq J(\square_k, p, q) - J(\square_n, p, q).
\]

**Proof.** For any $z \in Z_{n, k}$, since $v(\square_n) \in \mathcal{A}(z + \square_k)$, we use the quadratic response eq. (4.29) for $J(z + \square_k, p, q)$ that

\[
\mathbb{E}_\rho \left[ \frac{1}{\rho|\square|} \int_{z+\square} \frac{1}{2} |\nabla v(\square_n) - \nabla v(z + \square_k)|^2 \, d\mu \right] \leq \mathbb{E}_\rho \left[ \frac{1}{\rho|\square|} \int_{z+\square} \frac{1}{2} (\nabla v(\square_n) - \nabla v(z + \square_k)) \cdot a(\nabla v(\square_n) - \nabla v(z + \square_k)) \, d\mu \right] - J(z + \square_k, p, q) \]

\[
= J(z + \square_k, p, q) - \mathbb{E}_\rho \left[ \frac{1}{\rho|\square|} \int_{z+\square} \left( -\frac{1}{2} \nabla v(\square_n) \cdot a \nabla v(\square_n) - p \cdot a \nabla v(\square_n) + q \cdot \nabla v(\square_n) \right) \, d\mu \right].
\]
We sum this expression over all \( z \in Z_{n,k} \) to obtain that
\[
\frac{1}{|Z_{n,k}|} \sum_{z \in Z_{n,k}} \mathbb{E}_\rho \left[ \frac{1}{\rho[z_k]} \int_{z+\Box_k} \frac{1}{2} |\nabla v(\Box_n) - \nabla v(z + \Box_k)|^2 \, d\mu \right]
\leq \frac{1}{|Z_{n,k}|} \sum_{z \in Z_{n,k}} \left( J(z + \Box_k, p, q) - \frac{1}{\rho[z_k]} \int_{z+\Box_k} \left( \frac{1}{2} \nabla v(\Box_n) \cdot a \nabla v(\Box_n) - p \cdot a \nabla v(\Box_n) + q \cdot \nabla v(\Box_n) \right) \, d\mu \right)
= J(\Box_k, p, q) - J(\Box_n, p, q).
\]

In the last step, we use the stationarity of \( J \) and also eq. (4.20) for \( v(\Box_n) \).
\[\square\]

5. Quantitative Rate of Convergence

We are now ready to prove Theorem 2.1. We decompose the argument into a series of four steps.

5.1. Step 1: setup. We use the shorthand \( \mathbf{a}_n := \mathbf{a}_*(\Box_n) \), so that by eq. (4.28), the average slope of the function \( v(\cdot, \Box_n, p, q) \) is \( \mathbf{a}_n^{-1}q - p \), in the sense that
\[
\mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_n]} \int_U \nabla v(\cdot, \Box_n, p, q) \, d\mu \right] = \mathbf{a}_n^{-1}q - p.
\]

We let \( \tau_n \) denote a measure of the defect in the subadditivity of \( J \), precisely,
\[
\tau_n := \sup_{p, q \in B_1} (J(\Box_n, p, q) - J(\Box_{n+1}, p, q))
\]
\[
= \sup_{p \in B_1} (v(\Box_n, p) - v(\Box_{n+1}, p)) + \sup_{q \in B_1} (v^*(\Box_n, q) - v^*(\Box_{n+1}, q)).
\]

A direct corollary from eq. (5.2) is that for any integers \( n < m \),
\[
|\mathbf{a}_m^{-1} - \mathbf{a}_n^{-1}| = \sup_{q \in B_1} q \cdot (\mathbf{a}_m^{-1} - \mathbf{a}_n^{-1})q = \sup_{q \in B_1} (v^*(\Box_n, q) - v^*(\Box_m, q)) \leq C \sum_{k=n}^{m-1} \tau_k.
\]

We recall that \( \{\mathbf{a}(\Box_m)\}_{m \geq 0} \) is decreasing and \( \{\mathbf{a}_*(\Box_m)\}_{m \geq 0} \) is increasing, with the comparison \( \mathbf{a}_*(\Box_m) \leq \mathbf{a}(\Box_m) \). From eq. (4.21), we know that
\[
|\mathbf{a}(\Box_m) - \mathbf{a}| \leq |\mathbf{a}(\Box_m) - \mathbf{a}_*(\Box_m)| \leq C \sup_{p \in B_1} (J(\Box_m, p, \mathbf{a}_m p))^{1/2}.
\]

From now on, we thus fix \( p \in B_1 \), and focus on estimating \( J(\Box_m, p, \mathbf{a}_m p) \). We also assume without further notification that \( m \) is sufficiently large that \( 3^m \geq R_0 \), for the constant \( R_0 \) appearing in Proposition 3.9. We use \( A_{3^m+2}v(\cdot, \Box_{m+1}, p, \mathbf{a}_m p) \) to compare with eq. (4.27) and apply the quadratic response eq. (4.29). In the rest of Step 1, we write \( v(U) \) as a shorthand for \( v(\cdot, U, p, \mathbf{a}_m p) \), and decompose
\[
(J(\Box_m, p, \mathbf{a}_m p))^{1/2} = \left( \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_m]} \int_{\Box_m} \frac{1}{2} \nabla v(\Box_m) \cdot a \nabla v(\Box_m) \, d\mu \right] \right)^{1/2}
\leq \text{eq. (5.4)-a} + \text{eq. (5.4)-b},
\]
with
\[
\text{eq. (5.4)-a}
= \left( \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_m]} \int_{\Box_m} \frac{1}{2} (\nabla v(\Box_m) - \nabla A_{3^m+2}v(\Box_{m+1})) \cdot a(\nabla v(\Box_m) - \nabla A_{3^m+2}v(\Box_{m+1})) \, d\mu \right] \right)^{1/2},
\]
\[
\text{eq. (5.4)-b}
= \left( \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_m]} \int_{\Box_m} \frac{1}{2} \nabla v(\Box_m) \cdot a \nabla v(\Box_m) \, d\mu \right] \right)^{1/2}.
\]
We treat the two terms separately. For eq. (5.4)-a, since $A_{3m+2}^m$ is $A(m)$ (see Proposition A.2 for details), we use eq. (4.29) to get

$$|\text{eq. (5.4)-a}|^2$$

$$= J(\square_m, p, \bar{a}_m p) - \mathbb{E}_p \left[ \frac{1}{\rho(\square_m)} \int_{\square_m} \left( -\frac{1}{2} \nabla v(\square_{m+1}) \cdot \bar{a}_m \nabla v(\square_{m+1}) \right) \, d\mu \right]$$

$$= \mathbb{E}_p \left[ \int_{\square_m} (p \cdot a \nabla v(\square_{m+1}) + \bar{a}_m p \cdot \nabla A_{3m+2}^m(\square_{m+1})) \, d\mu \right].$$

Using Jensen’s inequality, we have

$$\mathbb{E}_p \left[ \int_{\square_m} \left( \frac{1}{2} \nabla A_{3m+2}^m(\square_{m+1}) \cdot a \nabla A_{3m+2}^m(\square_{m+1}) \right) \, d\mu \right] \leq \mathbb{E}_p \left[ \int_{\square_m} \left( \frac{1}{2} \nabla v(\square_{m+1}) \cdot a \nabla v(\square_{m+1}) \right) \, d\mu \right],$$

and the conditional expectation also implies that

$$\mathbb{E}_p \left[ \int_{\square_m} (p \cdot a \nabla v(\square_{m+1}) + \bar{a}_m p \cdot \nabla v(\square_{m+1})) \, d\mu \right]$$

$$= \mathbb{E}_p \left[ \int_{\square_m} (p \cdot a \nabla v(\square_{m+1}) + \bar{a}_m p \cdot \nabla v(\square_{m+1})) \, d\mu \right].$$

Thus we combine these terms with the quadratic response eq. (4.29) to obtain

$$|\text{eq. (5.4)-a}|^2 \leq J(\square_m, p, \bar{a}_m p) - J(\square_{m+1}, p, \bar{a}_m p) \leq C(d, \Lambda) \tau_m,$$

where the quantity $\tau_m$ is defined in eq. (5.2).

For the term eq. (5.4)-b, we can apply the modified Caccioppoli inequality eq. (3.37): there exist two finite positive constants $C(d, \Lambda)$ and $\theta(d, \Lambda)$ such that

$$\mathbb{E}_p \left[ \frac{1}{\rho(\square_m)} \int_{\square_m} \nabla A_{3m+2}^m(\square_{m+1}) \cdot a \nabla A_{3m+2}^m(\square_{m+1}) \, d\mu \right]$$

$$\leq \frac{C}{32m^2 \rho(\square_{m+1})} \mathbb{E}_p \left[ (v(\square_{m+1}))^2 \right] + \theta \mathbb{E}_p \left[ \frac{1}{\rho(\square_{m+1})} \int_{\square_{m+1}} \nabla v(\square_{m+1}) \cdot a \nabla v(\square_{m+1}) \, d\mu \right].$$

Using eq. (4.27), we see that the averaged gradient term on the right side of eq. (5.6) is $J(\square_{m+1}, p, \bar{a}_m p)$, and eq. (4.30) asserts that $J(\square_{m+1}, p, \bar{a}_m p) \leq J(\square_m, p, \bar{a}_m p)$. Therefore, we get the bound for eq. (5.4)-b

$$|\text{eq. (5.4)-b}|^2 \leq \frac{C}{32m^2 \rho(\square_{m+1})} \mathbb{E}_p \left[ (v(\square_{m+1}))^2 \right] + \theta J(\square_m, p, \bar{a}_m p).$$
We put eq. (5.5) and eq. (5.7) back to eq. (5.4), obtaining

\[ (J(\square_m, p, \bar{a}_m p))^{\frac{1}{2}} \leq C \tau_m^{\frac{1}{2}} + \left( \frac{C}{3\beta m \rho(\square_{m+1})} \right)^{\frac{1}{2}} \left[ \|v(\square_{m+1})\|^2_{L^2} + \theta J(\square_m, p, \bar{a}_m p) \right]^{\frac{1}{2}} \]

\[ \leq C \tau_m^{\frac{1}{2}} + \left( \frac{C}{3m \rho(\square_{m+1})} \right)^{\frac{1}{2}} \left[ \|v(\square_{m+1})\|^2_{L^2} + \theta J(\square_m, p, \bar{a}_m p) \right]^{\frac{1}{2}}. \]

Since \( \theta < 1 \), this gives

\[ J(\square_m, p, \bar{a}_m p) \leq C \left( \tau_m + \frac{C}{3^2 \rho(\square_{m+1})} \|v(\mu, \square_{m+1}, p, \bar{a}_m p)\|^2_{L^2} \right). \]

5.2. Step 2: flatness estimate. In this step, we estimate the \( L^2 \)-flatness of optimizers of \( J \). Notice that, using the result of Lemma 5.10 with \( v(\cdot, \square_{m+1}, p, \bar{a}_m p) \), the corresponding affine function is 0 and we obtain from eq. (5.8) that

\[ J(\square_m, p, \bar{a}_m p) \leq C \left( 3^{-\beta m} + \sum_{n=0}^{m} 3^{-\beta(m-n)} \tau_n \right). \]

**Lemma 5.1** (\( L^2 \)-flatness estimate). There exist \( \beta(d) > 0 \) and \( C(d, \Lambda, \rho) < \infty \) such that for every \( p, q \in B_1 \) and \( m \in \mathbb{N} \),

\[ \frac{1}{\rho(\square_{m+1})} \|v(\cdot, \square_m, p, q) - \ell_{\pi_{m+1}^{-1} q-p, \square_{m+1}}\|^2_{L^2} \leq C 3^{2m} \left( 3^{-\beta m} + \sum_{n=0}^{m} 3^{-\beta(m-n)} \tau_n \right). \]

**Proof.** In the rest of the proof, we write \( v(U) := v(\cdot, U, p, q) \) as we will not change \( p, q \) in the proof. Since \( E \left[ v(\square_{m+1}) - \ell_{\pi_{m+1}^{-1} q-p, \square_{m+1}} | \mathcal{G}_{m+1} \right] = 0 \), we can use the multiscale Poincaré inequality eq. (3.23)

\[ \frac{1}{\rho(\square_{m+1})} \|v(\square_{m+1}) - \ell_{\pi_{m+1}^{-1} q-p, \square_{m+1}}\|^2_{L^2} \]

\[ \leq C \left( E \left[ \frac{1}{\rho(\square_{m+1})} \int_{\square_{m+1}} |\nabla v(\square_{m+1}) - (\pi_{m+1}^{-1} q-p)|^2 \mathrm{d} \mu \right] \right)^{\frac{1}{2}} \]

\[ + C \sum_{n=0}^{m+1} 3^{n} \left( E \left[ \frac{1}{\rho(\square_{m+1})} \int_{\square_{m+1}} |S_{m+1} n \nabla v(\square_{m+1}) - (\pi_{m+1}^{-1} q-p)|^2 \mathrm{d} \mu \right] \right)^{\frac{1}{2}}. \]

The first term on the right side above is of constant order, by eq. (4.27). For the second term, we use a two-scale comparison for every \( 0 \leq n \leq m+1 \) that

\[ \left( E \left[ \frac{1}{\rho(\square_{m+1})} \int_{\square_{m+1}} |S_{m+1} n \nabla v(\square_{m+1}) - (\pi_{m+1}^{-1} q-p)|^2 \mathrm{d} \mu \right] \right)^{\frac{1}{2}} \]

\[ \leq |\pi_{m+1}^{-1} - \bar{a}_n^{-1}| + \left( E \left[ \frac{1}{\rho(\square_{m+1})} \sum_{z \in Z_{m+1} n} \int_{z+\square_{n}} |S_{m+1} n \nabla v(z+\square_{n}) - S_{m+1} n \nabla v(z+\square_{n})|^2 \mathrm{d} \mu \right] \right)^{\frac{1}{2}} \]

For the first term \( |\pi_{m+1}^{-1} - \bar{a}_n^{-1}| \) we have

\[ |\pi_{m+1}^{-1} - \bar{a}_n^{-1}|^2 \leq C(d, \Lambda) |\pi_{m+1}^{-1} - \bar{a}_n^{-1}|^2 \leq \sum_{k=n}^{m-1} \tau_n. \]
For the second term in eq. (5.12), recalling eq. (3.15), we use Jensen’s inequality and eq. (4.31) to get
\[
\mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \sum_{z \in \mathbb{Z}_{m+1,n}} \int_{z+\Box_n} \left| S_{m+1,n} \nabla v(\Box_{m+1}) - S_{m+1,n} \nabla v(z + \Box_n) \right|^2 d\mu \right] \\
\leq \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \sum_{z \in \mathbb{Z}_{m+1,n}} \int_{z+\Box_n} \left| \nabla v(\Box_{m+1}) - \nabla v(z + \Box_n) \right|^2 d\mu \right] \\
\leq \sum_{k=n}^{m} \tau_k.
\]
For the third term eq. (5.12), we use eq. (3.17), Jensen’s inequality, and stationarity. Here we remark that the operator $S_{n,n}^\ast$ is a conditional expectation with more information than $S_{m+1,n}$.
\[
\mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \sum_{z \in \mathbb{Z}_{m+1,n}} \int_{z+\Box_n} \left| S_{m+1,n} \nabla v(z + \Box_n) - (\mathbf{a}_n^{-1} q - p) \right|^2 d\mu \right] \\
\leq \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \sum_{z \in \mathbb{Z}_{m+1,n}} \int_{z+\Box_n} \left| S_{n,n}^\ast \nabla v(z + \Box_n) - (\mathbf{a}_n^{-1} q - p) \right|^2 d\mu \right] \\
= \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \int_{\Box_n} \left| S_n \nabla v(\Box_n) - (\mathbf{a}_n^{-1} q - p) \right|^2 d\mu \right].
\]
The estimation of this term is postponed to the next step. We will prove in Lemma 5.2 below that
\[
\mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \int_{\Box_n} \left| S_n \nabla v(\Box_n) - (\mathbf{a}_n^{-1} q - p) \right|^2 d\mu \right] \leq C3^{-\beta n} + \sum_{k=0}^{n-1} 3^{-\beta(n-k)} \tau_k.
\]
We put these estimates back to eq. (5.11) and obtain that
\[
\frac{1}{(\rho[\Box_{m+1}])^\frac{1}{2}} \left\| v(\Box_{m+1}) - \ell \mathbf{a}_n^{-1} q - p, \Box_{m+1} \right\|_{\mathcal{F}^2} \leq C \sum_{n=0}^{m} 3^n \left( 3^{-\beta n} + \sum_{k=0}^{n-1} 3^{-\beta(n-k)} \tau_k + \sum_{k=n}^{m} \tau_k \right)^\frac{1}{2}.
\]
We square the two sides and use the Cauchy-Schwarz inequality to obtain
\[
\frac{1}{(\rho[\Box_{m+1}])^\frac{1}{2}} \left\| v(\Box_{m+1}) - \ell \mathbf{a}_n^{-1} q - p, \Box_{m+1} \right\|_{\mathcal{F}^2}^2 \\
\leq C \left( \sum_{n=0}^{m} 3^n \right) \left( \sum_{n=0}^{m} 3^n \left( 3^{-\beta n} + \sum_{k=0}^{n-1} 3^{-\beta(n-k)} \tau_k + \sum_{k=n}^{m} \tau_k \right) \right) \\
\leq C3^{2m} \left( 3^{-\beta n} + \sum_{n=0}^{m} 3^{-\beta(m-n)} \tau_n \right),
\]
as announced. □

5.3. Step 3: variance estimate. In this part, we prove the following variance estimate, which was used in Step 2.

**Lemma 5.2** (Variance estimate). There exist $\beta(d) > 0$ and $C(d, \Lambda, \rho) < \infty$ such that for every $p, q \in B_1$ and $n \in \mathbb{N}$,
\[
(5.13) \quad \mathbb{E}_\rho \left[ \frac{1}{\rho[\Box_{m+1}]} \int_{\Box_n} \left| S_n \nabla v(\mu, \Box_n, p, q) - (\mathbf{a}_n^{-1} q - p) \right|^2 d\mu \right] \leq C3^{-\beta n} + \sum_{k=0}^{n-1} 3^{-\beta(n-k)} \tau_k.
\]
We denote by (5.15), (5.14) and (5.17) for approximate eq. (5.13) in scale 3

Proof. In the rest of the proof, we write \(v(U) := v(\cdot, U, p, q)\), as we will not change \(p, q\) in the proof. From eq. (5.1), we know that the average slope of \(v(\square_n)\) is \((\bar{a}_n^{-1} q - p)\), and notice that \(v(\square_n)\) is \(\mathcal{F}_{B_1(\square_n)}\)-measurable. Thus the idea is to use \(\{v(z + \square_k)\}_{z \in \mathbb{Z}^d_{n,k}}\) to approximate eq. (5.13) in scale 3\(^k\) with some error, and then apply the independence for \(v(z + \square_k)\) and \(v(z' + \square_k)\) for \(dist(z, z')\) large. However, different from the standard elliptic setting, here we will see a renormalization with random weights.

We start by relaxing eq. (5.13) to \(\mathcal{G}_{n,n-2}\). We observe that in fact \(S_n \nabla v(\square_n)\) is constant in \(\square_n\), so

\[
\int_{\square_n} |S_n \nabla v(\square_n) - (\bar{a}_n^{-1} q - p)|^2 d\mu = \frac{1}{\mu(\square_n)} \int_{\square_n} (S_n \nabla v(\square_n) - (\bar{a}_n^{-1} q - p))^2 d\mu
\]

We denote by \(V_n\) the left hand side of eq. (5.13). By triangle inequality, we have

\[
(V_n)^{\frac{1}{2}} \leq \text{eq. (5.14)-a} + \text{eq. (5.14)-b} + \text{eq. (5.14)-c},
\]

with

\[
\text{eq. (5.14)-a} = |\bar{a}_n^{-1} - \bar{a}_{n-2}^{-1}|
\]

\[
\text{eq. (5.14)-b} = \left( E_p \left[ \frac{1}{\rho(\square_n)} \sum_{z \in \mathbb{Z}^d_{n,n-2}} (S_n \nabla v(\square_n) - S_n \nabla v(z + \square_{n-2}))^2 d\mu \right]^\frac{1}{2} \right)^2
\]

\[
\text{eq. (5.14)-c} = \left( E_p \left[ \frac{1}{\rho(\square_n)} \sum_{z \in \mathbb{Z}^d_{n,n-2}} (S_n \nabla v(z + \square_{n-2}) - (\bar{a}_{n-2}^{-1} q - p))^2 d\mu \right]^\frac{1}{2} \right)^2
\]

The term eq. (5.14)-a can be controlled by eq. (5.3):

\[
\text{eq. (5.14)-a} \leq C(\tau_{n-2} + \tau_{n-1})^{\frac{1}{2}}.
\]

For the term eq. (5.14)-b, recalling eq. (3.17) and eq. (3.15), we use Jensen’s inequality and the two-scale comparison eq. (4.31) to get

\[
\text{eq. (5.14)-b} \leq \left( E_p \left[ \frac{1}{\rho(\square_n)} \sum_{z \in \mathbb{Z}^d_{n,n-2}} \int_{\square_n} |S_n \nabla v(\square_n) - S_n \nabla v(z + \square_{n-2})|^2 d\mu \right]^\frac{1}{2} \right)^2
\]

\[
\leq \left( E_p \left[ \frac{1}{\rho(\square_n)} \sum_{z \in \mathbb{Z}^d_{n,n-2}} \int_{\square_n} |\nabla v(\square_n) - \nabla v(z + \square_{n-2})|^2 d\mu \right]^\frac{1}{2} \right)^2
\]

\[
\leq (\tau_{n-2} + \tau_{n-1})^{\frac{1}{2}}.
\]

The term eq. (5.14)-c is the key for our result. To simplify a little more the notation, we write

\[
\begin{align*}
X_z := S_n \nabla v(z + \square_{n-2})(\mu, z) - (\bar{a}_{n-2}^{-1} q - p), \\
m_z := \mu (z + \square_{n-2}).
\end{align*}
\]

Notice that \(X_z, m_z\) are \(\mathcal{F}_{z + \square_{n-2}}\)-measurable. With this notation in place, we have

\[
\int_{z + \square_{n-2}} (S_n \nabla v(z + \square_{n-2}) - (\bar{a}_{n-2}^{-1} q - p)) d\mu = m_z X_z,
\]

and by eq. (5.1),

\[
E_p[m_z X_z] = 0.
\]
The term eq. (5.14)-c we want to estimate can be rewritten as

\[ \text{eq. (5.14)-c} = \left( \mathbb{E}_\rho \left[ \frac{1}{\rho|\mathbb{D}|} \left( \sum_{z \in \mathbb{Z}_{n-2}} \frac{m_z X_z}{m_z} \right)^2 \right] \right)^{\frac{1}{2}}. \]

If the coefficients \( m_z \) were deterministic, then we would be able to leverage on the finite range of dependence of \( X_z \) in this variance term. However, since the number of particles \( m_z \) is random, we introduce the event

\[ \mathcal{C}_{n,\rho,\delta} := \left\{ \mu \in \mathcal{M}_d(\mathbb{R}^d) : \forall z \in \mathbb{Z}_{n-2}, \left| \frac{\mu(z + \mathbb{D}_{n-2})}{\rho|\mathbb{D}_{n-2}|} - 1 \right| \leq \delta, \text{ and } \left| \frac{\mu(\mathbb{D}_n)}{\rho|\mathbb{D}_n|} - 1 \right| \leq \delta \}, \]

thus we can divide eq. (5.14)-c into two terms

\[ \text{eq. (5.14)-c} \leq \text{eq. (5.14)-c1} + \text{eq. (5.14)-c2}, \]

\[ \text{eq. (5.14)-c1} = \left( \mathbb{E}_\rho \left[ 1_{\{\mathcal{C}_{n,\rho,\delta}\}^c} \left( \frac{\sum_{z \in \mathbb{Z}_{n-2}} m_z X_z}{\sum_{z \in \mathbb{Z}_{n-2}} m_z} \right)^2 \right] \right)^{\frac{1}{2}}, \]

\[ \text{eq. (5.14)-c2} = \left( \mathbb{E}_\rho \left[ 1_{\{\mathcal{C}_{n,\rho,\delta}\}} \left( \frac{\sum_{z \in \mathbb{Z}_{n-2}} m_z X_z}{\sum_{z \in \mathbb{Z}_{n-2}} m_z} \right)^2 \right] \right)^{\frac{1}{2}}. \]

For the term eq. (5.14)-c1, we know that \( (\mathcal{C}_{n,\rho,\delta})^c \) is not typical in large scales, and we have the Chernoff bound

\[ \mathbb{P}_\rho[\mu \notin \mathcal{C}_{n,\rho,\delta}] \leq 3^{2d+1} \exp \left( -\frac{\rho|\mathbb{D}_{n-2}|\delta^2}{4} \right). \]

Moreover, by the Cauchy-Schwarz inequality,

\[ \frac{\left( \sum_{z \in \mathbb{Z}_{n-2}} m_z X_z \right)^2}{\sum_{z \in \mathbb{Z}_{n-2}} m_z} \leq \sum_{z \in \mathbb{Z}_{n-2}} m_z |X_z|^2. \]

We need a bound for the term \( |X_z|^2 \): recalling the definition in eq. (3.15) and eq. (4.28),

\[ S_{n-2,n-2}^z \nabla v(z + \mathbb{D}_{n-2})(\mu, z) = \mathbb{E}_\rho \left[ \int_{z + \mathbb{D}_{n-2}} \nabla v(z + \mathbb{D}_{n-2}) \, d\mu \ \bigg| \mathcal{G}_{n-2,n-2} \right] \]

\[ = a(z + \mathbb{D}_{n-2}; \mathcal{G}_{z + \mathbb{D}_{n-2}})^{-1} q - p. \]

Using the martingale structure of eq. (3.17), we have

\[ X_z = S_{n-2,n-2} \nabla v(z + \mathbb{D}_{n-2})(\mu, z) - (\overline{a}_{n-2} q - p) \]

\[ = \mathbb{E}_\rho \left[ \int_{z + \mathbb{D}_{n-2}} S_{n-2,n-2}^z \nabla v(z + \mathbb{D}_{n-2}) \, d\mu \ \bigg| \mathcal{G}_{n-2,n-2} \right] - (\overline{a}_{n-2}^{-1} q - p) \]

\[ = \mathbb{E}_\rho \left[ a(z + \mathbb{D}_{n-2}; \mathcal{G}_{z + \mathbb{D}_{n-2}})^{-1} - \overline{a}_{n-2}^{-1} q \right. \bigg| \mathcal{G}_{n-2,n-2} \bigg] q. \]

Then we use Jensen’s inequality and the bound of \( L \in a(z + \mathbb{D}_{n-2}; \mathcal{G}_{z + \mathbb{D}_{n-2}}) \leq \Lambda d \)

\[ |X_z|^2 = |S_{n-2,n-2} \nabla v(z + \mathbb{D}_{n-2}) - (\overline{a}_{n-2}^{-1} q - p)|^2 \]

\[ = \mathbb{E}_\rho \left[ a(z + \mathbb{D}_{n-2}; \mathcal{G}_{z + \mathbb{D}_{n-2}})^{-1} \bigg| \mathcal{G}_{n-2,n-2} \right] \leq \Lambda^2. \]
This concludes that

\begin{equation}
\text{eq. (5.14)-c1} \leq A^2 \mathbb{E}_\rho \left[ \frac{1 \{ \mathcal{C}_{n,\rho,\delta} \} }{\rho(\square_n)} \mu(\square_n) \right] \leq C(d, A) \frac{1}{\rho(\square_n)} \exp \left( -\frac{\rho(\square_n-\delta^2)}{4} \right) \\
\leq C(d, A, \rho) 3^{-d_n}.
\end{equation}

Finally, we treat eq. (5.14)-c2. We calculate eq. (5.14)-c2 at first with the conditional expectation with respect to \( \mathcal{G}_{n,n-2} \). Clearly, \( \mathcal{C}_{n,\rho,\delta} \) is \( \mathcal{G}_{n,n-2} \)-measurable, and under this condition \( \mu(\square_n) \geq (1 - \delta) \rho(\square_n) \), so we have

\begin{equation}
|\text{eq. (5.14)-c2}|^2 = \frac{1}{\rho(\square_n)} \mathbb{E}_\rho \left[ \frac{1 \{ \mathcal{C}_{n,\rho,\delta} \} }{\mu(\square_n)} \mathbb{E}_\rho \left[ \left( \sum_{z \in \mathbb{Z}^{n,n-2}} m_z X_z \right)^2 \mid \mathcal{G}_{n,n-2} \right] \right] \\
\leq \frac{1}{\rho(\square_n)} \mathbb{E}_\rho \left[ \frac{1}{(1 - \delta) \rho(\square_n)} \mathbb{E}_\rho \left[ 1 \{ \mathcal{C}_{n,\rho,\delta} \} \left( \sum_{z \in \mathbb{Z}^{n,n-2}} m_z X_z \right)^2 \mid \mathcal{G}_{n,n-2} \right] \right] .
\end{equation}

We would like to develop the term \( |\sum_{z \in \mathbb{Z}^{n,n-2}} m_z X_z|^2 \) and also drop out the indicator term. The argument here is deterministic

\[
\left| \sum_{z \in \mathbb{Z}^{n,n-2}} m_z X_z \right|^2 = \sum_{z,z' \in \mathbb{Z}^{n,n-2}} m_z m_{z'} X_z \cdot X_{z'} + \sum_{z,z' \in \mathbb{Z}^{n,n-2}} |m_z m_{z'} X_z \cdot X_{z'}| \\
\leq \frac{1}{2} \sum_{z,z' \in \mathbb{Z}^{n,n-2}} (m_z)^2 |X_z|^2 + (m_{z'})^2 |X_{z'}|^2 + \sum_{z,z' \in \mathbb{Z}^{n,n-2}} m_z m_{z'} X_z \cdot X_{z'},
\]

where \( |z - z'|_\infty := \max_{1 \leq i \leq d} |z_i - z'_i| \). We now add back the indicator \( 1 \{ \mathcal{C}_{n,\rho,\delta} \} \) and develop it

\begin{equation}
1 \{ \mathcal{C}_{n,\rho,\delta} \} \left| \sum_{z \in \mathbb{Z}^{n,n-2}} m_z X_z \right|^2 \\
\leq 1 \{ \mathcal{C}_{n,\rho,\delta} \} \left( (1 + \delta) \rho(\square_n) \sum_{z,z' \in \mathbb{Z}^{n,n-2}} (m_z |X_z|^2 + m_{z'} |X_{z'}|^2) + \sum_{z,z' \in \mathbb{Z}^{n,n-2}} m_z m_{z'} X_z \cdot X_{z'} \right) \\
\leq \frac{(1 + \delta) \rho(\square_n)}{2} \sum_{z,z' \in \mathbb{Z}^{n,n-2}} (m_z |X_z|^2 + m_{z'} |X_{z'}|^2) + \sum_{z,z' \in \mathbb{Z}^{n,n-2}} m_z m_{z'} X_z \cdot X_{z'} .
\end{equation}

From the first line to the second line above, we use that \( m_z \leq (1 + \delta) \rho(\square_n) \) under the event \( \mathcal{C}_{n,\rho,\delta} \). We then keep in mind that the quantity in \( \cdots \) on the second line of eq. (5.22) is always larger than \( |\sum_{z \in \mathbb{Z}^{n,n-2}} m_z X_z|^2 \), so it is nonnegative. Therefore, from the second line to the third line, we can drop the indicator function in front.
Inserting this estimate into eq. (5.21), we obtain that
\[
|\text{eq. (5.14)}-c2|^2 \leq \frac{1}{\rho|\square_n|} \left( \frac{1 + \delta}{|\square_n|} \right) \sum_{z,z' \in \mathbb{Z}_{n-2}} \mathbb{E}_\rho \left[ \frac{1}{2} \left( |m_zX_z|^2 + m_z|X_{z'}|^2 \right) \right] + \frac{1}{\rho|\square_n| (1 - \delta) |\square_n|} \sum_{z,z' \in \mathbb{Z}_{n-2}} \mathbb{E}_\rho \left[ m_zm_zX_z \cdot X_{z'} \right].
\]

The sum in the second line is 0, because for $|z - z'|_\infty \geq 3^{n-1}$, $m_zX_z$ and $m_zX_{z'}$ are independent,
\[
\mathbb{E}_\rho \left[ m_zm_zX_z \cdot X_{z'} \right] = \mathbb{E}_\rho \left[ m_zX_z \right] \cdot \mathbb{E}_\rho \left[ m_zX_{z'} \right] = 0.
\]
For the sum in the first line, $\mathbb{E}_\rho [m_z|X_z|^2]$ is nothing but
\[
\mathbb{E}_\rho \left[ \int_{z + \square_n} |S_{n,n-2} \nabla v(z + \square_n) - (\tilde{\alpha}_{n-2} q - p) |^2 \ d\mu \right],
\]
We use Jensen’s inequality to shrink the operator to $S_{n-2,n-2}^\ast$ that
\[
\mathbb{E}_\rho \left[ \int_{z + \square_n} |S_{n,n-2} \nabla v(z + \square_n) - (\tilde{\alpha}_{n-2} q - p) |^2 \ d\mu \right] \leq \mathbb{E}_\rho \left[ \int_{z + \square_n} |S_{n-2,n-2}^\ast \nabla v(z + \square_n) - (\tilde{\alpha}_{n-2} q - p) |^2 \ d\mu \right] = \mathbb{E}_\rho \left[ \int_{\square_n} |S_{n-2} \nabla v(\square_n) - (\tilde{\alpha}_{n-2} q - p) |^2 \ d\mu \right].
\]
There are at most $9^d \times 5^d$ pairs $z, z' \in \mathbb{Z}_{n-2}$ such that $|z - z'|_\infty < 3^{n-1}$; see Figure 3 for an illustration. Therefore, we obtain
\[
|\text{eq. (5.14)}-c2|^2 \leq \frac{5}{9}^d \left( \frac{1 + \delta}{1 - \delta} \right) \mathbb{E}_\rho \left[ \frac{1}{\rho|\square_n|} \int_{\square_n} |S_{n-2} \nabla v(\square_n) - (\tilde{\alpha}_{n-2} q - p) |^2 \ d\mu \right] = \frac{5}{9}^d \left( \frac{1 + \delta}{1 - \delta} \right) V_{n-2},
\]
where we recall that $V_n$ is the left hand side of eq. (5.13). We put this estimate together with eq. (5.15), (5.16), (5.20) back to eq. (5.14) to obtain the recurrence relation
\[
(V_n)^\frac{1}{2} \leq \frac{5}{9}^d \left( \frac{1 + \delta}{1 - \delta} \right)^\frac{1}{2} (V_{n-2})^{\frac{1}{2}} + C(\tau_{n-2} + \tau_{n-1})^{\frac{3}{2}} + C3^{-dn}.
\]
By choosing $\delta(d) > 0$ sufficiently small, we obtain the desired result eq. (5.13). □

5.4. Step 4: iterations. Once we obtain the estimate eq. (5.9), it remains to do some numerical iterations, similarly to [11, Page 59-60]. For the reader’s convenience, we recall the main steps here. Let $\{e_i\}_{1 \leq i \leq d}$ denote the canonical basis in $\mathbb{R}^d$, and define
\[
F_m := \sum_{i=1}^d J(\square_m, e_i, \tilde{\alpha}_m e_i).
\]
In order to obtain an exponential decay for $(F_m)_{m \geq 0}$, we first introduce a weighted version of this quantity:
\[
\tilde{F}_m := \sum_{n=0}^m 3^{-\frac{\delta}{2}(m-n)} F_n.
\]
we have
\begin{equation}
(5.23)
\end{equation}
following, we calculate some bounds for $q_{I/d}$ using also eq. (4.19), that equation of type $\beta$
Here the exponent $\beta$ is the same as in eq. (5.9). It is clear that $F_m \leq \tilde{F}_m$, so it suffices
to prove an exponential decay for $(\tilde{F}_m)_{m \geq 0}$. We will do so by proving a recurrence
equation of type $\tilde{F}_{m+1} \leq C(\tilde{F}_m - \tilde{F}_{m+1})$ for some constant $C(d, \Lambda) < \infty$. Thus in the following, we calculate some bounds for $(\tilde{F}_m - \tilde{F}_{m+1})$ and $\tilde{F}_{m+1}$.

Starting with $(\tilde{F}_m - \tilde{F}_{m+1})$, we write
\[ \tilde{F}_m - \tilde{F}_{m+1} \geq \sum_{n=0}^{n} 3^{-\beta (n-1)} (F_n - F_{n+1}) - C3^{-\beta n}\]
Noticing that $\bar{a}_{n+1}p$ is the minimizer for the mapping $q \mapsto J(\square_{n+1} + p, q)$ in eq. (4.19),
we have
\begin{equation}
(5.23)
F_{n+1} = \sum_{i=1}^{d} J(\square_{n+1}, e_i, \bar{a}_{n+1}e_i) \leq \sum_{i=1}^{d} J(\square_{n+1}, e_i, \bar{a}_{n}e_i).
\end{equation}
Using also eq. (4.19), that $\Lambda d \leq a_{n} \leq \Lambda d$, and that $p \mapsto \nu(\square_{n}, p) - \nu(\square_{n+1}, p)$ and
$q \mapsto \nu^{*}(\square_{n}, q) - \nu^{*}(\square_{n+1}, q)$ are positive semidefinite quadratic forms, we get
\[
F_n - F_{n+1} \geq \sum_{i=1}^{d} (J(\square_{n}, e_i, \bar{a}_{n}e_i) - J(\square_{n+1}, e_i, \bar{a}_{n}e_i))
\]
\[
= \sum_{i=1}^{d} (\nu(\square_{n}, e_i) - \nu(\square_{n+1}, e_i)) + \sum_{i=1}^{d} (\nu^{*}(\square_{n}, a_{n}e_i) - \nu^{*}(\square_{n+1}, a_{n}e_i))
\]
\[
\geq C^{-1} \left( \sup_{p \in B_{1}} (\nu(\square_{n}, p) - \nu(\square_{n+1}, p)) + \sup_{q \in B_{1}} (\nu^{*}(\square_{n}, q) - \nu^{*}(\square_{n+1}, q)) \right)
\]
\[
\geq C^{-1} \tau_{n},
\]
Figure 3. In the cube $\square_{n}$ and all its sub-cubes $\{z + \square_{n-2}\}_{z \in \square_{n,n-2}}$, for a chosen sub-cube $z_{0} + \square_{n-2}$ (the cube in dark red), the support of $v(z_{0} + \square_{n-2})$ is in $z_{0} + \square_{n-1}$ (the cube in light red), so it has at most $5^d$ cubes of scale $3^{n-2}$ whose associated function has a support intersecting with $z_{1} + \square_{n-1}$ (the cube in blue). For example, $v(z_{2} + \square_{n-2})$ has correlation with $v(z_{0} + \square_{n-2})$, while $v(z_{1} + \square_{n-2})v(z_{3} + \square_{n-2})$ do not. This gives us the contraction factor $(\frac{5}{9})^d$. 

Quantitative Homogenization of Interacting Particle Systems
and thus

\[(5.24) \quad F_{m} - F_{m+1} \geq C^{-1} \sum_{n=0}^{m} 3^{-\frac{\beta}{2}(m-n)} \tau_n - C 3^{-\frac{\beta m}{2}}.\]

For the upper bound of \(F_{m+1}\), we use eq. (5.23) to see that \(F_n \leq F_{n+1}\), so

\[\tilde{F}_{m+1} = 3^{-\frac{\beta}{2}(m+1)} F_0 + \sum_{n=0}^{m} 3^{-\frac{\beta}{2}(m-n)} F_{n+1}\]

\[\leq C 3^{-\frac{\beta m}{2}} + \sum_{n=0}^{m} 3^{-\frac{\beta}{2}(m-n)} F_n.\]

Then we apply eq. (5.9) into the result above to get

\[(5.25) \quad \tilde{F}_{m+1} \leq C 3^{-\frac{\beta m}{2}} + \sum_{n=0}^{m} 3^{-\frac{\beta}{2}(m-n)} \left(3^{-\beta \delta n} + \sum_{k=0}^{m} 3^{-\beta (n-k)} \tau_k \right)\]

\[\leq C 3^{-\frac{\beta m}{2}} + C 3^{-\frac{\beta m}{2}} \sum_{k=0}^{m} 3^{-\frac{\beta}{2}(2k-n)} \tau_k.\]

We combine eq. (5.24) and eq. (5.25), to obtain \(C(\tilde{F}_{m} - \tilde{F}_{m+1} + C 3^{-\frac{\beta m}{2}}) \geq \tilde{F}_{m+1}\), which implies

\[\tilde{F}_{m+1} \leq \theta \tilde{F}_{m} + C 3^{-\frac{\beta m}{2}},\]

for some \(\theta(d, \Lambda) \in (0, 1)\). We thus conclude for the exponential decay of \((\tilde{F}_m)_{m \geq 0}\), and thus also of \(F_m\), since \(F_m \leq \tilde{F}_m\). By eq. (4.21), this completes the proof of Theorem 2.1.

**Appendix A. Some elementary properties of the function spaces**

**Lemma A.1** (Canonical projection). Let \(f : \mathcal{M}_d(\mathbb{R}^d) \to \mathbb{R}\) be a function, and for every Borel set \(U\), measure \(\mu \in \mathcal{M}_d(\mathbb{R}^d)\), and \(n \in \mathbb{N}\), let \(f_n(. \mu \sqcup U^c)\) denote the (permutation-invariant) function

\[f_n(. \mu \sqcup U^c) : \left\{ \begin{array}{ll} U^n & \to \mathbb{R} \\ (x_1, \ldots, x_n) & \mapsto f(\sum_{i=1}^{n} \delta_{x_i} + \mu \sqcup U^c). \end{array} \right.\]

The following statements are equivalent.

1. The function \(f\) is \(\mathcal{F}\)-measurable.
2. For every \(n \in \mathbb{N}\), the function \(f_n\) is \(\mathcal{B}^{\otimes n}_U \otimes \mathcal{F}_{U^c}\)-measurable.

**Proof.** We start from (1) \(\Rightarrow\) (2). Because \(\mathcal{F} = \mathcal{F}_U \otimes \mathcal{F}_{U^c}\), it suffices to study the product function

\[f = 1_{(\mu(V_1) = n_1)} 1_{(\mu(V_2) = n_2)} 1_{(\mu(U) = n)}\]

for some Borel sets \(V_1 \subseteq U, V_2 \subseteq U^c\). In this case, we have

\[\{f_n = 1\} = \{\mu(V_1) = n_1\} \cap \{\mu(V_2) = n_2\} \cap \{\mu(U) = n\}\]

\[= \bigcup_{\sigma \in S_n} \left( \bigcap_{i=1}^{n_1} \{x_{\sigma(i)} \in V_1\} \bigcap_{j=n_1+1}^{n} \{x_{\sigma(j)} \in (U \setminus V_1)\} \bigcap \{\mu(V_2) = n_2\} \right),\]

where \(S_n\) is the symmetric group. This proves that \(f_n\) is \(\mathcal{B}^{\otimes n}_U \otimes \mathcal{F}_{U^c}\)-measurable.
We turn to (2) ⇒ (1). Let us pick a suitable $f_n$ and $\mu_L U = \sum_{i=1}^n \delta_{x_i}$, then the main point is to establish the $\mathcal{F}$-measurable property. Since $f_n$ is $\mathcal{B}^n \otimes \mathcal{F}_{\nu}$-measurable and permutation-invariant, it suffices to study the function of type

\begin{equation} \tag{A.1} f_n = \sum_{\sigma \in S_n} \left( \prod_{i=1}^n 1_{\{x_{\sigma(i)} \in V_i\}} \right) 1_{\{\mu U^c(V_0) = n_0\}} 1_{\{\mu(U) = n\}}, \end{equation}

for $\{V_i\}_{0 \leq i \leq n}$ Borel sets. This is still a complicated function, but we can add one more condition

\begin{equation} \tag{A.2} \forall 1 \leq i, j \leq n, \quad V_i = V_j \text{ or } V_i \cap V_j = \emptyset. \end{equation}

For example, let $\{\tilde{V}_j\}_{0 \leq j \leq n}$ be all the different elements in $\{V_i\}_{0 \leq i \leq n}$, and $\tilde{V}_j$ appears $n_j$ times. For the functions of type eq. (A.1) satisfying the condition eq. (A.2), the $\mathcal{F}$-measurable property is easy to treat since we have

\[ \sum_{\sigma \in S_n} \left( \prod_{i=1}^n 1_{\{x_{\sigma(i)} \in V_i\}} \right) 1_{\{\mu U^c(V_0) = n_0\}} 1_{\{\mu(U) = n\}} = \left( \prod_{j=1}^m 1_{\{\mu(V_j) = n_j\}} \right) 1_{\{\mu U^c(V_0) = n_0\}} 1_{\{\mu(U) = n\}}, \]

which is an $\mathcal{F}$-measurable function.

Finally, let us conclude that for a general $f_n$ in eq. (A.1), they can be decomposed into the sum of the functions with the propriety eq. (A.2). Let us see the case $n = 2$, where we have the following decomposition

\[ 1_{\{x_1 \in V_1\}} 1_{\{x_2 \in V_2\}} = \left( 1_{\{x_1 \in (V_1 \cap V_2)\}} + 1_{\{x_1 \in (V_1 \setminus V_2)\}} \right) \left( 1_{\{x_2 \in (V_2 \cap V_1)\}} + 1_{\{x_2 \in (V_2 \setminus V_1)\}} \right) \]

\[ = 1_{\{x_1 \in (V_1 \cap V_2)\}} 1_{\{x_2 \in (V_2 \cap V_1)\}} + 1_{\{x_1 \in (V_1 \setminus V_2)\}} 1_{\{x_2 \in (V_2 \setminus V_1)\}} + 1_{\{x_1 \in (V_1 \setminus V_2)\}} 1_{\{x_2 \in (V_2 \cap V_1)\}} + 1_{\{x_1 \in (V_1 \setminus V_2)\}} 1_{\{x_2 \in (V_2 \setminus V_1)\}}. \]

For a general $n$, one can use induction and this concludes the proof. \hfill \Box

**Proposition A.2.** For every $s > 0$ and $f \in \mathcal{H}^1(Q_s)$, we have $A_s f \in \mathcal{H}^1(Q_s)$, and for every $x \in \text{supp}(\mu) \cap Q_s$,

\begin{equation} \tag{A.3} \nabla (A_s f)(\mu, x) = A_s (\nabla f)(\mu, x). \end{equation}

Moreover, if $s > 2$ and $f \in \mathcal{A}(Q_s)$, then $A_s f \in \mathcal{A}(Q_{s-2})$.

**Proof.** At first, we should remark the well-definedness of the right side of eq. (A.3). Notice that the Poisson measure can be decomposed as a sum of the independent parts $\mu = \mu \mathcal{L} Q_s + \mu^c \mathcal{L} Q_s^c$, we have

\[ A_s f = \int_{M_s(\mathbb{R}^d)} f(\mu \mathcal{L} Q_s + \mu^c \mathcal{L} Q_s^c) \, d\mathbb{P}_\rho(\mu'). \]

Thus the right-hand side of eq. (A.3) is defined as

\begin{equation} \tag{A.4} A_s (\nabla f)(\mu, x) := \int_{M_s(\mathbb{R}^d)} \nabla f(\mu \mathcal{L} Q_s + \mu^c \mathcal{L} Q_s^c, x) \, d\mathbb{P}_\rho(\mu'). \end{equation}

We prove eq. (A.3) and $A_s f \in \mathcal{H}^1(Q_s)$ for the functions in $\mathcal{C}^\infty(Q_s) \cap \mathcal{H}^1(Q_s)$ as they are dense, and we can focus on the case $\mu(Q_s) = n$ fixed. We use Lemma A.1 to write

\[ f \left( \sum_{i=1}^n \delta_{x_i} + \mu \mathcal{L} Q_s^c \right) = f_n(x_1, \cdots, x_n, \mu \mathcal{L} Q_s^c). \]
Following the property of product measure, for every \((x_1, x_2, \ldots, x_n) \in (Q_s)^n\), the mapping

\[
\mu \llcorner Q_s \mapsto f_n(x_1, \ldots, x_n, \mu \llcorner Q_s),
\]

is \(\mathcal{F}_{Q_s}\)-measurable. Thus for every \((x_1, x_2, \ldots, x_n) \in (Q_s)^n\), the mapping

\[
\mu \llcorner Q_s^c \mapsto \nabla f_n(x_1, \ldots, x_n, \mu \llcorner Q_s^c),
\]

is also \(\mathcal{F}_{Q_s}\)-measurable because it is the limit of \(\mathcal{F}_{Q_s}\)-measurable functions. Then we observe that

\[
\|\nabla f_n\|_{L^\infty((Q_s)^n)} = \sup_{(Q_s)^n} \left( \sum_{k=1}^n \left| \nabla f_n(x_1, \ldots, x_n, \mu \llcorner Q_s^c)^2 \right| \right)^{\frac{1}{2}} = \sup_{(Q_s)^n} \left( \sum_{k=1}^n \left| \nabla f_n(x_1, \ldots, x_n, \mu \llcorner Q_s^c)^2 \right| \right)^{\frac{1}{2}},
\]

as a supremum of a countable number of \(\mathcal{F}_{Q_s}\)-measurable functions, is finite and \(\mathcal{F}_{Q_s}\)-measurable. Thus we can define a cut-off version of \(f\) that

\[
f^{n,M} = f1_{\{\mu(Q_s)=n\}}1_{\{\|\nabla f_n\|_{L^\infty((Q_s)^n)} \leq M\}},
\]

and we can establish eq. (A.3) at first for \(f^{n,M}\). For every \(x \in Q_s \cap \text{supp}(\mu)\), we have

\[
\partial_k(A_s f^{n,M})(\mu, x) = \lim_{h \to 0} \int_{\mathcal{M}_d(\mathbb{R}^d)} f((\mu - \delta_x + \delta_{x+he_k}) \llcorner Q_s + \mu' \llcorner Q_s^c) - f((\mu \llcorner Q_s + \mu' \llcorner Q_s^c)) \frac{dP}{\rho'(\mu')1_{\{\mu(Q_s)=n\}}} d\mu(x + \theta e_k),
\]

for \(h\) small enough such that \(x + he_k \in Q_s\). Since \(f \in C^\infty(Q_s)\), we use Lemma A.1 and the mean value theorem

\[
\frac{f(\mu - \delta_x + \delta_{x+he_k}) - f(\mu)}{h} = \partial_k f(\mu - \delta_x + \delta_{x+he_k}, x + \theta e_k),
\]

for some \(\theta \in (0, 1)\). With the indicator \(1_{\{\|\nabla f_n\|_{L^\infty((Q_s)^n)} \leq M\}}\), this term is bounded by \(M\), so we can use the dominated convergence theorem that

\[
\partial_k(A_s f^{n,M})(\mu, x) = \int_{\mathcal{M}_d(\mathbb{R}^d)} \lim_{h \to 0} \frac{f((\mu - \delta_x + \delta_{x+he_k}) \llcorner Q_s + \mu' \llcorner Q_s^c) - f((\mu \llcorner Q_s + \mu' \llcorner Q_s^c))}{h} \frac{dP}{\rho'(\mu')1_{\{\mu(Q_s)=n\}}} d\mu(x + \theta e_k) = \int_{\mathcal{M}_d(\mathbb{R}^d)} \partial_k f(\mu \llcorner Q_s + \mu' \llcorner Q_s^c) \frac{dP}{\rho'(\mu')} \frac{1}{\mathcal{M}_d(\mathbb{R}^d)} d\mu(Q_s) \frac{dP}{\rho'(\mu')}1_{\{\mu(Q_s)=n\}},
\]
which implies that\(A_\varepsilon f_\varepsilon^{n,M} \in \mathcal{H}^1(Q_s)\). Then we use once again Jensen’s inequality for \(f^{n,M}\) and \(f^{n,M'}\) with \(M < M'\)

\[
\begin{align*}
\mathbb{E}_\rho \left[ \int_{Q_s} |\nabla (A_\varepsilon f_\varepsilon^{n,M})|^2 (\mu, x) \, d\mu(x) \right]
&= \mathbb{E}_\rho \left[ \int_{Q_s} \left| \int_{M_\varepsilon(\mathbb{R}^d)} \nabla f^{n,M}(\mu L \overline{Q}_s + \mu' L \overline{Q}_s^c, x) \, d\mathbb{P}_\rho(\mu') \right|^2 \, d\mu(x) \right] \\
&\leq \mathbb{E}_\rho \left[ \int_{Q_s} \left| \nabla f^{n,M} - \nabla f^{n,M'} \right|^2 (\mu, x) \, d\mu(x) \right] \\
&= \mathbb{E}_\rho \left[ \int_{Q_s} \left| \nabla f \right|^2 (\mu, x) \, d\mu(x) \right].
\end{align*}
\]

So \(\{f^{n,M}\}_{M \geq 1}\) gives a Cauchy sequence in \(\mathcal{H}^1(Q_s)\), and the only candidate is \(f \mathbb{1}_{\{\mu(\overline{Q}_s) = n\}}\) because it is the limit in \(L^2\). By this and a linear combination, we establish eq. (A.3) for \(f \in C^\infty(Q_s) \cap \mathcal{H}^1(Q_s)\), and we can then extend to a general function in \(\mathcal{H}^1(Q_s)\) by the density argument.

For the part of \(a\)-harmonic function, we suppose \(f \in \mathcal{A}(Q_s)\) and test \(\phi \in \mathcal{H}^1_0(Q_{s-2})\) with eq. (A.3),

\[
\begin{align*}
\mathbb{E}_\rho \left[ \int_{Q_{s-2}} (\nabla A_\varepsilon f)(\mu, x) \cdot a(\mu, x) \nabla \phi(\mu, x) \, d\mu(x) \right]
&= \mathbb{E}_\rho \left[ \int_{Q_{s-2}} A_\varepsilon(\nabla f)(\mu, x) \cdot a(\mu, x) \nabla \phi(\mu, x) \, d\mu(x) \right] \\
&= \mathbb{E}_\rho \left[ \int_{Q_{s-2}} \left( \int_{M_\varepsilon(\mathbb{R}^d)} \nabla f(\mu L \overline{Q}_s + \mu' L \overline{Q}_s^c, x) \, d\mathbb{P}_\rho(\mu') \right) \cdot a(\mu, x) \nabla \phi(\mu, x) \, d\mu(x) \right].
\end{align*}
\]

Restricted on \(x \in Q_{s-2}\), we have \(a(\mu, x), \nabla \phi(\mu, x)\) are \(\mathcal{F}_{Q_s} \otimes \mathcal{B}_{Q_s^{-1}}\)-measurable, so we have

\[
\forall x \in \text{supp}(\mu) \cap Q_{s-2}, \quad a(\mu, x) \nabla \phi(\mu, x) = a(\mu L \overline{Q}_s, x) \nabla \phi(\mu L \overline{Q}_s, x).
\]

We can enter the part in the integration, and then use Fubini’s lemma

\[
\begin{align*}
\mathbb{E}_\rho \left[ \int_{Q_{s-2}} (\nabla A_\varepsilon f)(\mu, x) \cdot a(\mu, x) \nabla \phi(\mu, x) \, d\mu(x) \right]
&= \mathbb{E}_\rho \left[ \int_{Q_{s-2}} \left( \int_{M_\varepsilon(\mathbb{R}^d)} \nabla f(\mu L \overline{Q}_s + \mu' L \overline{Q}_s^c, x) \cdot a(\mu L \overline{Q}_s, x) \nabla \phi(\mu L \overline{Q}_s, x) \, d\mathbb{P}_\rho(\mu') \right) \, d\mu(x) \right] \\
&= \mathbb{E}_\rho \left[ \int_{Q_{s-2}} \nabla f(\mu, x) \cdot a(\mu, x) \nabla \phi(\mu, x) \, d\mu(x) \right] \\
&= 0.
\end{align*}
\]

In the last step, we use \(f \in \mathcal{A}(Q_s)\) and this finishes the proof. \(\square\)
APPENDIX B. EQUIVALENT DEFINITIONS OF THE EFFECTIVE DIFFUSION MATRIX

Recall that we defined $\bar{a}(U)$ and $\bar{a}_*(U)$ according to eq. (4.1)-eq. (4.3). The proof of Theorem 2.1 ensures the existence of a constant $C < \infty$ and an exponent $\alpha > 0$ such that for every $m \in \mathbb{N}$,

$$|\bar{a}(\square_m) - \bar{a}| + |\bar{a}_*(\square_m) - \bar{a}| \leq C 3^{-\alpha m}. \tag{B.1}$$

Throughout this appendix, we will only rely on the qualitative statement that

$$\bar{a} = \lim_{m \to \infty} \bar{a}(\square_m) = \lim_{m \to \infty} \bar{a}_*(\square_m). \tag{B.2}$$

The first main goal of this appendix is to demonstrate that the definition we chose for the bulk diffusion matrix indeed coincides with the “stationary” definition appearing in works such as [69, 34]. Adapted to our context, this alternative definition takes the following form. For

$$\Gamma := \left\{ \mu \mapsto \int_{\mathbb{R}^d} \tau_x g(\mu) \, dx, \quad g \in C_c^\infty(\mathbb{R}^d) \cap H_0^1(\mathbb{R}^d) \right\}, \tag{B.3}$$

we let $\bar{a}$ be the $d$-by-$d$ matrix such that for every $p \in \mathbb{R}^d$,

$$p \cdot \bar{a}p := \inf_{u \in \Gamma} \mathbb{E}_p \left[ \left( p + \nabla u(\mu + \delta_0, 0) \right) \cdot \bar{a}(\mu + \delta_0, 0) \left( p + \nabla u(\mu + \delta_0, 0) \right) \right], \tag{B.4}$$

where in eq. (B.3), we used the notation $\tau_x g(\mu) := g(\tau_x \mu)$. Notice that for $g \in C_c^\infty(\mathbb{R}^d) \cap H_0^1(\mathbb{R}^d)$ and $u : \mu \mapsto \int_{\mathbb{R}^d} \tau_x g(\mu) \, dx \in \Gamma$, the function $\mu \mapsto u(\mu)$ is typically not well-defined unless $\mu$ is of finite support. However, the quantity $\nabla u(\mu, \cdot)$ makes sense whenever the measure $\mu$ is $\sigma$-finite, since in the sum $\nabla u(\mu, y) = \int_{\mathbb{R}^d} \nabla(\tau_x g)(\mu, y) \, dx$, the function $g$ is local, and thus the integrand $\nabla(\tau_x g)(\mu, y)$ is non-zero only for $x$ in a bounded set. With this interpretation of $\nabla u$, the right side of eq. (B.4) is well-defined.

**Theorem B.1.** We have $\bar{a} = \bar{a}$.

The second main goal of this appendix is to demonstrate that the infimum in eq. (B.4) is achieved in a suitable completion of the space $\Gamma$. We also show that the optimizer, which we call the (stationary) corrector, can be obtained as a limit of approximations based on the finite-volume optimizers for $\nu$ or $\nu^*$. Denoting

$$\mathcal{M}_\delta(\mathbb{R}^d) := \left\{ (\mu, x) \in \mathcal{M}_\delta(\mathbb{R}^d) \times \mathbb{R}^d : x \in \text{supp} \mu \right\},$$

we introduce the space

$$\mathcal{L}^2 := \left\{ f : \mathcal{M}_\bullet(\mathbb{R}^d) \to \mathbb{R} : f \text{ is measurable and } \mathbb{E}_\nu \left[ \int_{\mathbb{R}^d} |f(\mu, x)|^2 \, d\mu(x) \right] < \infty \right\},$$

and its local version

$$\mathcal{L}^2_{\text{loc}} := \left\{ f : \mathcal{M}_\bullet(\mathbb{R}^d) \to \mathbb{R} : f \text{ is measurable and } \text{for every compact } K \subseteq \mathbb{R}^d, \quad \mathbb{E}_\nu \left[ \int_K |f(\mu, x)|^2 \, d\mu(x) \right] < \infty \right\}. \tag{B.5}$$

In these definitions, we say that $f : \mathcal{M}_\bullet(\mathbb{R}^d) \to \mathbb{R}$ is measurable provided that the mapping

$$\left\{ \begin{array}{c} \mathcal{M}_\delta(\mathbb{R}^d) \times \mathbb{R}^d \to \mathbb{R} \\ (\mu, x) \mapsto f(\mu, x)1_{\{x \in \text{supp} \mu\}} \end{array} \right\}$$
is $\mathcal{F} \otimes \mathcal{B}$-measurable. The space $\mathcal{L}^2_{\text{loc}}$ is naturally endowed with the family of seminorms

$$\begin{align*}
\mathcal{L}^2_{\text{loc}} & \to \mathbb{R} \\
 f & \mapsto \mathbb{E}_\phi \left[ \int_K |f(\mu, x)|^2 \, d\mu(x) \right]^\frac{1}{2},
\end{align*}$$

indexed by all the compact sets $K \subseteq \mathbb{R}^d$. This family of seminorms turns $\mathcal{L}^2_{\text{loc}}$ into a complete space.

For every $p \in \mathbb{R}^d$ and $m \in \mathbb{N}$, we let $\phi(\cdot, \square_m, p)$ be such that the minimizer in the definition of $\nu(\square_m, p)$ is $\ell_{p, \square_m} + \phi(\cdot, \square_m, p)$, where we recall that $\ell_{p, \square_m}$ was defined in eq. (B.4). Similarly, we let $\phi^*(\cdot, \square_m, p)$ be such that $\ell_{p, \square_m} + \phi^*(\cdot, \square_m, p)$ is the maximizer in the definition of $\nu^*(\square_m, \bar{a}_*(\square_m) p)$. More precisely, using the notation introduced in Proposition 4.1, we have

$$\nu(\cdot, \square_m, p) = \ell_{p, U} + \phi(\cdot, \square_m, p),$$

and

$$u(\cdot, \square_m, \bar{a}_*(\square_m) p) = \ell_{p, \square_m} + \phi^*(\cdot, \square_m, p).$$

Finally, we define $\nabla \tilde{\phi}_{p,m}$ according to the formula

$$(B.6) \quad \tilde{\phi}_{p,m} : \mu \mapsto \frac{1}{|\square_m|} \int_{\mathbb{R}^d} \tau_x \phi(\mu, \square_m, p) \, dx = \frac{1}{|\square_m|} \int_{\mathbb{R}^d} \phi(\mu, x + \square_m, p) \, dx.$$  

Notice that, while $\tilde{\phi}_{p,m}(\mu)$ is ill-defined when $\mu \sim \text{Poi}(\rho)$, the quantity $\nabla \tilde{\phi}_{p,m}(\mu, \cdot)$ is still well-defined, for the same reason as in the discussion following eq. (B.4). Our second main result is as follows.

**Theorem B.2.** The following statements hold for every $p \in \mathbb{R}^d$.

1. The sequence $(\nabla \tilde{\phi}_{p,m})_{m \in \mathbb{N}}$ is a Cauchy sequence in $(\mathcal{L}^2_{\text{loc}})^d$. Its limit, which we denote by $\nabla \tilde{\phi}_p$, satisfies

$$(B.7) \quad \forall v \in \mathcal{H}_0^1(\mathbb{R}^d), \quad \mathbb{E}_p \left[ \int_{\mathbb{R}^d} \nabla v \cdot a(p + \nabla \tilde{\phi}_p) \, d\mu \right] = 0.$$  

2. We have

$$(B.8) \quad \lim_{m \to \infty} \mathbb{E}_p \left[ \frac{1}{|\square_m|} \int_{\square_m} |\nabla \phi(\mu, \cdot, \square_m, p) - \nabla \phi_p(\mu, \cdot)|^2 \, d\mu \right] = 0,$$

as well as

$$(B.9) \quad \lim_{m \to \infty} \mathbb{E}_p \left[ \frac{1}{|\square_m|} \int_{\square_m} |\nabla \phi^*(\mu, \cdot, \square_m, p) - \nabla \phi_p(\mu, \cdot)|^2 \, d\mu \right] = 0.$$  

3. The effective diffusion matrix $\bar{a}$ satisfies

$$(B.10) \quad p \cdot \bar{a} p = \mathbb{E}_p \left[ (p + \nabla \phi_p(\mu + \delta_0, 0)) \cdot a(\mu + \delta_0, 0) (p + \nabla \phi_p(\mu + \delta_0, 0)) \right],$$

as well as

$$(B.11) \quad \bar{a} p = \mathbb{E}_p \left[ a(\mu + \delta_0, 0) (p + \nabla \phi_p(\mu + \delta_0, 0)) \right].$$

As a preparation towards the proof of these results, we state in the following proposition a number of elementary properties about the function space $\Gamma$.

**Proposition B.3.** Let $g \in \mathcal{C}_c^\infty(\mathbb{R}^d) \cap \mathcal{H}_0^1(\mathbb{R}^d)$ be an $\mathcal{F}_{\square_m}$-measurable function, and let $u := \int_{\mathbb{R}^d} \tau_y g \, dx \in \Gamma$. The following properties hold:

1. $u \mapsto \nabla u(\mu + \delta_y, y)$ is a stationary field, i.e. $\nabla u(\mu + \delta_y, y) = \nabla u(\tau_y \mu + \delta_0, 0)$.
2. $\nabla u$ has mean zero, that is,

$$(B.12) \quad \mathbb{E}_p \left[ \nabla u(\mu + \delta_0, 0) \right] = 0.$$
We put the definition of \( y \) to get
\[
\nabla \mu(k) = \frac{1}{\rho |\Box_n|} \int_{\Box_n} \nabla g(\mu, y) \cdot a(\mu, y) \nabla g(\mu, y) \, d\mu(y).
\]

**Proof.** (1) We use the definition to write
\[
\begin{align*}
\partial_h u(\mu + \delta, y) &= \lim_{h \to 0} \frac{1}{h} \left( \int_{\mathbb{R}^d} \tau_x g(\mu + \delta + \delta_x h) - \tau_x g(\mu + \delta) \, dx \right) \\
&= \lim_{h \to 0} \frac{1}{h} \left( \int_{y+\Box_n} \tau_x g(\mu + \delta + \delta_x h) - \tau_x g(\mu + \delta) \, dx \right) \\
&= \lim_{h \to 0} \frac{1}{h} \left( \int_{y+\Box_n} \tau_x g(\mu + \delta + \delta_x h) - \tau_x y g(\tau_y \mu + \delta_0 + h) \, dx \right).
\end{align*}
\]
From the first line to the second line, we used the fact that \( g \) is \( \mathcal{F}_{\Box_n} \)-measurable, so if the transport vector \( x \) does not belong to \( y + \Box_n \), then the integrand vanishes (up to a boundary layer that vanishes in the limit \( h \to 0 \)). We then do the change of variables \( z = x - y \) to get
\[
\partial_h u(\mu + \delta, y) = \lim_{h \to 0} \frac{1}{h} \left( \int_{\Box_n} \tau_z g(\tau_y \mu + \delta_0 + h) \, dx \right) = \partial_h u(\tau_y \mu + \delta_0, 0),
\]
which means that \( y \mapsto \nabla u(\mu + \delta, y) \) is a stationary gradient field.

(2) We use the equations developed in the last question. Since \( g \in \mathcal{C}_c^\infty(\mathbb{R}^d) \), we can exchange the integration and derivative and get
\[
\nabla u(\mu + \delta, y) = \int_{\mathbb{R}^d} \nabla \tau_z g(\tau_y \mu + \delta_0, 0) \, dz.
\]
We evaluate this gradient at \( y = 0 \)
\[
\begin{align*}
\mathbb{E}_\rho [\nabla u(\mu + \delta_0, 0)] &= \mathbb{E}_\rho \left[ \int_{\mathbb{R}^d} \nabla \tau_z g(\mu + \delta_0, 0) \, dz \right] \\
&= \mathbb{E}_\rho \left[ \int_{\mathbb{R}^d} \nabla g(\tau_z \mu + \delta, -z) \, dz \right] \\
&= \mathbb{E}_\rho \left[ \int_{\mathbb{R}^d} \nabla g(\mu + \delta, z) \, dz \right] \\
&= 0.
\end{align*}
\]
From the second line to the third line, we used the stationarity of the Poisson point process. Because \( g \in \mathcal{C}_c^\infty(\mathbb{R}^d) \), then \( \int_{\mathbb{R}^d} \nabla g(\mu + \delta, z) \, dz = 0 \) and the integration in the third line vanishes.

(3) We pick a cube \( Q_L = \left( -\frac{L}{2}, \frac{L}{2} \right)^d \) with \( L > 0 \), make use of the stationarity of \( y \mapsto \nabla u(\mu + \delta, y) \) and Mecke’s identity (see [52, Theorem 4.1])
\[
\mathbb{E}_\rho [\nabla u(\mu + \delta_0, 0) \cdot a(\mu + \delta_0, 0) \nabla u(\mu + \delta_0, 0)]
\]
\[
= \mathbb{E}_\rho \left[ \frac{1}{|Q_L|} \int_{Q_L} \nabla u(\mu + \delta, y) \cdot a(\mu + \delta, y) \nabla u(\mu + \delta, y) \, dy \right] \\
= \mathbb{E}_\rho \left[ \frac{1}{|Q_L|} \int_{Q_L} \nabla u(\mu, y) \cdot a(\mu, y) \nabla u(\mu, y) \, d\mu(y) \right].
\]
We put the definition \( \mu = \int_{\mathbb{R}^d} \tau_x g \, dx \) into the equation. For the gradient at \( y \), as \( g \) is \( \mathcal{F}_{\Box_n} \)-measurable, thus only the term \( \nabla(x-y)(\mu, y) \) for \( x \in y + \Box_n \) contributes. This
We next apply Jensen’s inequality to obtain that
\begin{align*}
B.15 \quad & \mathbb{E}_p \left[ \nabla u(\mu + \delta_0, 0) \cdot \mathbf{a}(\mu + \delta_0, 0) \nabla u(\mu + \delta_0, 0) \right] \\
& = \mathbb{E}_p \left[ \frac{1}{\rho |Q_L|} \int_{Q_L} \left( \int_{y + \vartriangle_n} \nabla (\tau_x g)(\mu, y) \, dx \right) \cdot \mathbf{a}(\mu, y) \left( \int_{y + \vartriangle_n} \nabla (\tau_x g)(\mu, y) \, dx \right) \, d\mu(y) \right].
\end{align*}

In the last line, we use Fubini’s lemma and exchange \( \int (\ldots) \, dx \) with \( \int (\ldots) \, d\mu(y) \). In this procedure, we have to enlarge the domain from \( Q_L \) to \( Q_{L+3^n} \), because for the gradient at \( y \in Q_L \), \( \nabla (\tau_x g)(\mu, y) \) contributes for the transport \( x \in Q_{L+3^n} \) (see Figure 4 as an illustration). Using the stationarity of the Poisson point process, we have
\begin{align*}
& \mathbb{E}_p \left[ \int_{y + \vartriangle_n} \nabla (\tau_x g)(\mu, y) : \mathbf{a}(\mu, y) \nabla (\tau_x g)(\mu, y) \, d\mu(y) \right] \\
& = \mathbb{E}_p \left[ \int_{\vartriangle_n} \nabla g(\mu, y) : \mathbf{a}(\mu, y) \nabla g(\mu, y) \, d\mu(y) \right].
\end{align*}

which helps us conclude that
\begin{align*}
& \mathbb{E}_p \left[ (\xi + \nabla u(\mu + \delta_0, 0)) : \mathbf{a}(\mu + \delta_0, 0)(\xi + \nabla u(\mu + \delta_0, 0)) \right] \\
& \leq 3^{2dn} |Q_{L+3^n}| \mathbb{E}_p \left[ \frac{1}{\rho |\vartriangle_n|} \int_{\vartriangle_n} \nabla g(\mu, y) : \mathbf{a}(\mu, y) \nabla g(\mu, y) \, d\mu(y) \right].
\end{align*}

We take \( L \to \infty \) and obtain the desired result. \( \square \)

**Remark B.4.** The inequality eq. (B.13) is essentially sharp when \( \nabla g \) itself is close to a stationary field. Indeed, if \( g \) is close to a stationary field, then
\[ \nabla \tau_x g(\mu, y) = \nabla g(\tau_x - \mu, y - x) \approx \nabla g(\mu, y), \]
which implies that the application of Jensen’s inequality in eq. (B.15) is essentially sharp. The error introduced by a boundary layer in a subsequent step of the proof disappears as we take \( L \to \infty \) at the end.

As a corollary of Proposition B.3, we can also propose the following equivalent definition of \( \tilde{\mathbf{a}} \).

**Corollary B.5.** For any open set \( U \subseteq \mathbb{R}^d \), we have
\begin{equation}
B.16 \quad \xi : \tilde{\mathbf{a}} \xi = \inf_{u \in U} \mathbb{E}_p \left[ \frac{1}{|U|} \int_U (\xi + \nabla u) : \mathbf{a}(\xi + \nabla u) \, d\mu \right].
\end{equation}

**Proof.** It is a direct result of Mecke’s identity (see [52, Theorem 4.1]) and the stationarity of \( y \mapsto \nabla u(\mu + \delta_y, y) \). \( \square \)

With the help of Proposition B.3, we can now prove the first main theorem of this appendix.
Proof of Theorem B.1. We decompose the proof into two steps.

Step 1: Bound from below $\tilde{a} \geq a$. We fix $m \in \mathbb{N}$ and a sequence of approximate minimizers $\{\phi_p^{(i)}\}_{i \geq 1}$ for the variational problem in eq. (B.4), which we write in the form

$$\phi_p^{(i)}(\mu) = \int_{\mathbb{R}^d} \tau_x g_i(\mu) \, dx$$

for some $g_i \in C_c^\infty(\mathbb{R}^d) \cap \mathcal{H}^1_0(\mathbb{R}^d)$. Now we propose a modified version in $\mathcal{H}^1(\square_m)$ defined by

$$\tilde{\phi}_p^{(i)}(\mu) := \int_{K_i} \tau_x g_i(\mu) \, dx,$$

with $K_i \subseteq \mathbb{R}^d$ a large compact set so that $\tilde{\phi}_p^{(i)} \in \mathcal{H}^1(\square_m)$ and

$$\forall y \in \square_m, \quad \nabla \tilde{\phi}_p^{(i)}(\mu, y) = \nabla \phi_p^{(i)}(\mu, y).$$

Then we test $p + \nabla \tilde{\phi}_p^{(i)}$ in the optimization problem for $\nu^*(\square_m, q)$ to get that

$$\frac{1}{2} q \cdot \tilde{a}^*_1(\square_m) q \geq \mathbb{E}_p \left[ \frac{1}{|\rho|_{\square_m}} \int_{\square_m} \left( -\frac{1}{2} (p + \nabla \tilde{\phi}_p^{(i)}) \cdot a(p + \nabla \tilde{\phi}_p^{(i)}) + q \cdot (p + \nabla \tilde{\phi}_p^{(i)}) \right) \, d\mu \right]$$

$$= \mathbb{E}_p \left[ \frac{1}{|\rho|_{\square_m}} \int_{\square_m} \left( -\frac{1}{2} (p + \nabla \phi_p^{(i)}) \cdot a(p + \nabla \phi_p^{(i)}) + q \cdot (p + \nabla \phi_p^{(i)}) \right) \, d\mu \right].$$

We use the stationarity of $y \mapsto \nabla \phi_p^{(i)}(\mu + \delta_y, y)$, eq. (B.12) and let $i \to \infty$ to obtain

$$\frac{1}{2} q \cdot \tilde{a}^*_1(\square_m) q \geq -\frac{1}{2} p \cdot \tilde{a} p + p \cdot q.$$

Taking $q = \tilde{a}_*(\square_m)p$ leads to

$$p \cdot \tilde{a} p \geq p \cdot a_*(\square_m)p.$$

Finally, we let $m \to \infty$ and conclude that $\tilde{a} \geq a$.

Step 2: Bound from above $\tilde{a} \leq a$. We hope to prove $\tilde{a} \leq a$ by testing the variational formula eq. (B.4) with a suitable candidate, namely the function $\tilde{\phi}_{p,m}$ introduced in eq. (B.6). Since $\phi(\cdot, \square_m, p) \in \mathcal{H}^1(\square_m)$, and using eq. (B.13), we can approximate
where the function validity of at least some of the statements in Theorem B.2. We now turn to the proof

\[ \nabla \tilde{\phi}_{p,m} \] in \((\mathcal{L}^2_{\text{loc}})^d\) arbitrarily closely with elements of \(\Gamma\). It thus follows that we can use \(\tilde{\phi}_{p,m}\) as a candidate in the variational problem in eq. (B.4), and use the comparison inequality eq. (B.13) to get that

\[ \begin{align*}
\langle \nabla \tilde{\phi}_{p,m}, \mu \rangle &\leq \mathbb{E}_\rho \left[ \int_{\mathbb{R}^d} \nabla \phi(x, y, z + \Box_n, p) \, dx \right] \\
&= p \cdot \tilde{a}(\Box_n) \rho.
\end{align*} \]

Finally, we let \(m \to \infty\) and conclude that \(\tilde{a} \leq \tilde{a}\).

Proof of part (1) of Theorem B.2. We decompose the proof into four steps.

Step 1: \(\{\nabla \tilde{\phi}_{p,m}\}_{m \geq 1}\) is a Cauchy sequence in \((\mathcal{L}^2_{\text{loc}})^d\). We fix \(n \leq m\) and, recalling the notation \(\mathcal{Z}_{m,n} := \mathbb{Z}^d \cap \Box_m\), we observe that

\[ \begin{align*}
\nabla \tilde{\phi}_{p,m,n}(\mu, y) &= \frac{1}{\Box_n} \int_{\mathbb{R}^d} \nabla \phi(x, y, x + \Box_n, p) \, dx \\
&= \frac{1}{\Box_m} \int_{\mathbb{R}^d} \sum_{z \in \mathcal{Z}_{m,n}} \nabla \phi(x, y, x + z + \Box_n, p) \, dx \\
&= \frac{1}{\Box_m} \int_{\mathbb{R}^d} \nabla \phi(x, y, x + \Box_n, p) \, dx,
\end{align*} \]

where the function \(\phi_{p,m,n}\) is defined as

\[ \phi_{p,m,n}(\mu) := \sum_{z \in \mathcal{Z}_{m,n}} \phi(z + \Box_n, p). \]

For any compact set \(K\), we use Mecke’s identity (see [52, Theorem 4.1]) and the stationarity of \(\nabla \tilde{\phi}_{p,m}\)

\[ \begin{align*}
\mathbb{E}_\rho \left[ \int_K |\nabla \tilde{\phi}_{p,m} - \nabla \tilde{\phi}_{p,n}|^2(\mu, y) \, d\mu(y) \right] \\
&= \rho \mathbb{E}_\rho \left[ \int_K |\nabla \tilde{\phi}_{p,m} - \nabla \tilde{\phi}_{p,n}|^2(\mu, y, \Box_n, p) \, d\mu(y) \right] \\
&= \rho |K| \mathbb{E}_\rho \left[ |\nabla \tilde{\phi}_{p,m} - \nabla \tilde{\phi}_{p,n}|^2(\mu, y, \Box_n, p) \right].
\end{align*} \]

Then we use the comparison inequality eq. (B.13) and obtain that

\[ \begin{align*}
\mathbb{E}_\rho \left[ \int_K |\nabla \tilde{\phi}_{p,m} - \nabla \tilde{\phi}_{p,n}|^2(\mu, y) \, d\mu(y) \right] \\
&\leq \rho |K| \mathbb{E}_\rho \left[ \int_{\Box_m} |\nabla \phi(x, y, x + \Box_n, p) - \nabla \phi_{p,m,n}(\mu, y)|^2 \, d\mu(y) \right] \\
&\leq \rho |K| \left( \nu(\Box_n, p) - \nu(\Box_m, p) \right).
\end{align*} \]

By eq. (B.2), this shows that \(\{\nabla \tilde{\phi}_{p,m}\}_{m \geq 1}\) is a Cauchy sequence in \((\mathcal{L}^2_{\text{loc}})^d\).

Step 2: Harmonic property - setting up. Denote the limit by \(\nabla \tilde{\phi}_{p}\), we set things up to prove the harmonic property eq. (B.7) by approximation. We fix \(s > 0\) and
\( v \in \mathcal{C}_c^\infty(\mathbb{R}^d) \cap \mathcal{K}_0^1(\mathbb{R}^d) \) which is \( \mathcal{F}_{Q_s} \)-measurable, and observe that

\[
\mathbb{E}_\rho \left[ \int_{\mathbb{R}^d} \nabla \cdot \mathbf{a}(p + \nabla \phi_p) \, d\mu \right] = \lim_{m \to \infty} \mathbb{E}_\rho \left[ \int_{Q_s} \nabla \cdot \mathbf{a}(p + \nabla \phi_{p,m}) \, d\mu \right] = \lim_{m \to \infty} \mathbb{E}_\rho \left[ \int_{Q_s} \nabla v(\mu, y) \cdot \mathbf{a}(\mu, y) \left( \int_{y + \Box_m} p + \nabla \phi(\mu, y, x + \Box_m, p) \, dx \right) \, d\mu(y) \right].
\]

We then use Fubini’s lemma to exchange the order of integration,

\[
\mathbb{E}_\rho \left[ \int_{\mathbb{R}^d} \nabla \cdot \mathbf{a}(p + \nabla \phi_p) \, d\mu \right] = \lim_{m \to \infty} \frac{1}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{Q_{3m-s}} \left( \int_{x + \Box_m} \nabla v(\mu, y) \cdot \mathbf{a}(\mu, y) \left( p + \nabla \phi(\mu, y, x + \Box_m, p) \right) \, d\mu(y) \right) \, dx \right].
\]

For \( m \) sufficiently large, we can decompose the domain of integration in \( x \) in the expression above into \( Q_{3m-s} \) and \( Q_{3m+s} \setminus Q_{3m-s} \). We analyse the contribution of each of these quantities in each of the following two steps.

**Step 3: Integration in \( Q_{3m-s} \).** Notice that for \( x \in Q_{3m-s} \), we have \( Q_s \subseteq x + \Box_m \) (see Figure 5 for an illustration), thus we can drop the indicator \( 1_{\{y \in Q_s\}} \) in the inner integral and use the \( \mathbf{a} \)-harmonic property of \( p + \nabla \phi(x + \Box_m, p) \) to get that

\[
\frac{1}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{Q_{3m-s}} \left( \int_{x + \Box_m} \nabla v(\mu, y) \cdot \mathbf{a}(\mu, y) \left( p + \nabla \phi(\mu, y, x + \Box_m, p) \right) \, d\mu(y) \right) \, dx \right] \]
\[
= \frac{1}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{x + \Box_m} \nabla v(\mu, y) \cdot \mathbf{a}(\mu, y) \left( p + \nabla \phi(\mu, y, x + \Box_m, p) \right) \, d\mu(y) \right] \]
\[
= 0.
\]

**Figure 5.** The red cube represents \( Q_s \), and the blue and green cubes represent respectively \( Q_{3m+s} \) and \( Q_{3m-s} \). For \( x \in Q_{3m-s} \), we have \( Q_s \subseteq x + \Box_m \); for \( x \notin Q_{3m+s} \), \( (x + \Box_m) \cap Q_s = \emptyset \); for \( x \in Q_{3m+s} \setminus Q_{3m-s} \), \( x + \Box_m \) and \( Q_s \) have non-empty intersection but \( Q_s \) is not totally contained in \( x + \Box_m \). These three cases are represented by \( z_1, z_2, z_3 \).
Step 4: Boundary layer $Q_{3^{m+s}} \setminus Q_{3^m - s}$. We use Young’s inequality to bound the term with the integral over $x \in Q_{3^{m+s}} \setminus Q_{3^m - s}$ by

$$
\begin{align*}
&\frac{\Lambda}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{Q_{3^{m+s}} \setminus Q_{3^m - s}} \left( \int_{(x+\Box_m) \cap Q_{3^m - s}} |\nabla v(\mu, y)|^2 + \left| p + \nabla \phi(\mu, y, x + \Box_m, p) \right|^2 \, d\mu(y) \right) \, dx \right] \\
\leq& \frac{\Lambda|Q_{3^{m+s}} \setminus Q_{3^m - s}|}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{Q_{3^m - s}} |\nabla \phi(\mu, y)|^2 \, d\mu(y) \right] \\
+& \frac{\Lambda}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{Q_{3^{m+s}} \setminus Q_{3^m - s}} \left( \int_{(x+\Box_m) \cap Q_{3^m - s}} \left| p + \nabla \phi(\mu, y, x + \Box_m, p) \right|^2 \, d\mu(y) \right) \, dx \right].
\end{align*}
$$

For the term (A), we have, for a constant $C$ that may depend on $s$,

$$(A) \leq C \Lambda 3^{-m} \mathbb{E}_\rho \left[ \int_{Q_{3^m - s}} |\nabla v(\mu, y)|^2 \, d\mu(y) \right] \xrightarrow{m \to \infty} 0.$$ 

For the term (B), we use the stationarity to observe that

$$
\mathbb{E}_\rho \left[ \int_{(x+\Box_m) \cap Q_{3^m - s}} \left| p + \nabla \phi(\mu, y, x + \Box_m, p) \right|^2 \, d\mu(y) \right] = \mathbb{E}_\rho \left[ \int_{\Box_m \cap (-x+Q_{3^m})} \left| p + \nabla \phi(\mu, y, \Box_m, p) \right|^2 \, d\mu(y) \right].
$$

We apply once again Fubini’s lemma to (B) and get that

$$(B) = \frac{\Lambda}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{Q_{3^{m+s}} \setminus Q_{3^m - s}} \left( \int_{\Box_m} \left| p + \nabla \phi(\mu, y, \Box_m, p) \right|^2 1_{\{y \in (-x+Q_{3^m})\}} \, d\mu(y) \right) \, dx \right] \\
= \frac{\Lambda}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{\Box_m} \left| p + \nabla \phi(\mu, y, \Box_m, p) \right|^2 \left( \int_{Q_{3^{m+s}} \setminus Q_{3^m - s}} 1_{\{y \in (-x+Q_{3^m})\}} \, d\mu(y) \right) \, dx \right] \\
\leq \frac{\Lambda \Lambda_Q}{|\Box_m|} \mathbb{E}_\rho \left[ \int_{\Box_m} \left| p + \nabla \phi(\mu, y, \Box_m, p) \right|^2 1_{\{\text{dist}(y, \partial \Box_m) \leq s\}} \, d\mu(y) \right].
$$

That this term converges to zero is a consequence of the stronger estimate given by Lemma B.6 below. This concludes the proof for $v \in \mathcal{C}^\infty_c(\mathbb{R}^d) \cap \mathcal{H}_0^1(\mathbb{R}^d)$, and then we can use density argument to extend to general $v \in \mathcal{H}_0^1(\mathbb{R}^d)$. \hfill \Box

In the proof above, we appealed to the following boundary layer estimate, which we state as a separate lemma for future reference (and which is stronger than what was needed for the purpose of the proof above, since the boundary layer size is allowed to increase with $m$).

**Lemma B.6** (Boundary layer estimate). For every sequence $(s_m)_{m \in \mathbb{N}}$ such that $s_m \leq 3^m$ and $\lim_{m \to \infty} 3^{-m} s_m = 0$, we have

$$
\lim_{m \to \infty} \mathbb{E}_\rho \left[ \frac{1}{|\Box_m|} \int_{\Box_m} \left| p + \nabla \phi(\mu, y, \Box_m, p) \right|^2 1_{\{\text{dist}(y, \partial \Box_m) \leq s_m\}} \, d\mu(y) \right] = 0.
$$

**Proof.** The idea is to make use of the renormalization argument. We define a mesoscopic scale $n$ associated to $m$ such that $s_m \leq 3^n$, $n \to \infty$ and $m - n \to \infty$. Then
we immediately have
\[
\mathbb{E}_p \left[ \frac{1}{\rho} \int_{\square_m} |p + \nabla \phi(\mu, y, \square_m, p)|^2 1_{\{\text{dist}(y, \partial \square_m) \leq \varepsilon_m\}} d\mu(y) \right] \\
\leq \mathbb{E}_p \left[ \frac{1}{\rho} \int_{\square_m} |p + \nabla \phi(\mu, y, \square_m, p)|^2 1_{\{\text{dist}(y, \partial \square_m) \leq 3^n\}} d\mu(y) \right] .
\]

We propose to compare \(\phi(\cdot, \square_m, p)\) with \(\phi_{p,m,n} \in H^1_0(\square_m)\) defined as in eq. (B.18):

\[
\phi_{p,m,n}(\mu) = \sum_{z \in \mathbb{Z}_{m,n}} \phi(\mu, z + \square_n, p) .
\]

Then we have
\[\begin{align*}
\mathbb{E}_p \left[ \frac{1}{\rho} \int_{\square_m} |p + \nabla \phi(\mu, y, \square_m, p)|^2 1_{\{\text{dist}(y, \partial \square_m) \leq 3^n\}} d\mu(y) \right] \\
\leq 2 \mathbb{E}_p \left[ \frac{1}{\rho} \int_{Q_{3m} \setminus Q_{3m-2,3^n}} |p + \nabla \phi_{p,m,n}(\mu, y)|^2 d\mu(y) \right] \\
+ 2 \mathbb{E}_p \left[ \frac{1}{\rho} \int_{Q_{3m} \setminus Q_{3m-2,3^n}} |\nabla \phi_{p,m,n}(\mu, y) - \nabla \phi(\mu, y, \square_m, p)|^2 d\mu(y) \right] .
\end{align*}\] (B.20)

For the first term eq. (B.20)-a, we do partition of sum into cubes of size \(3^n\), then we have
\[
\mathbb{E}_p \left[ \frac{1}{\rho} \int_{Q_{3m} \setminus Q_{3m-2,3^n}} |p + \nabla \phi_{p,m,n}(\mu, y)|^2 d\mu(y) \right] \\
= \frac{|\square_n|}{|\square_m|} \sum_{z \in \mathbb{Z}_{m,n} \cap (Q_{3m} \setminus Q_{3m-2,3^n})} \mathbb{E}_p \left[ \frac{1}{\rho} \int_{z + \square_n} |p + \nabla \phi(\mu, y, z + \square_n, p)|^2 d\mu(y) \right] \\
= \frac{|Q_{3m} \setminus Q_{3m-2,3^n}|}{|Q_{3m}|} \nu(\square_n, p) \\
\leq 3^{-(m-n)}\Lambda |\rho| .
\]

For the second term eq. (B.20)-b, we have
\[
\mathbb{E}_p \left[ \frac{1}{\rho} \int_{Q_{3m} \setminus Q_{3m-2,3^n}} |\nabla \phi_{p,m,n}(\mu, y) - \nabla \phi(\mu, y, \square_m, p)|^2 d\mu(y) \right] \\
\leq \mathbb{E}_p \left[ \frac{1}{\rho} \int_{Q_{3m}} |\nabla \phi_{p,m,n}(\mu, y) - \nabla \phi(\mu, y, \square_m, p)|^2 d\mu(y) \right] \\
= \nu(\square_n, p) - \nu(\square_m, p) .
\]

Therefore, when we take \(m, n \to \infty\), both eq. (B.20)-a and eq. (B.20)-b go to 0, so the boundary layer in a mesoscopic scale can be neglected. □

Now that the gradient of the whole-space corrector \(\nabla \phi_p\) is well-defined, we can proceed to complete the proof of Theorem B.2.

**Proof of parts (2) and (3) of Theorem B.2.** We start by discussing the validity of the identities eq. (B.10) and eq. (B.11). We use the stationary approximate corrector \(\phi_{p,n}\)
defined in eq. (B.6), and observe from eq. (B.17) and Theorem B.1 that

\[ p \cdot \bar{a} = \lim_{m \to \infty} \mathbb{E}_p \left[ \frac{1}{|\rho|} \left( p + \nabla \tilde{\phi}_{p,n}(\mu, \cdot, \square, p) \right) \cdot \mathbf{a}(\mu, \cdot, \square, p) \right] \]

The identity eq. (B.10) then follows from the convergence of \( \nabla \tilde{\phi}_{p,n} \) to \( \nabla \phi \) in \( L^2_{\mu,\text{loc}} \). For the second identity, we can use the fact that

\[ \mathbb{E}_p \left[ \frac{1}{|\rho|} \int_{\square_m} |\nabla \phi^*(\mu, \cdot, m, p) - \nabla \phi^*(\mu, \cdot)|^2 \, d\mu \right] \]

the estimate eq. (B.8), and the stationarity of \( \nabla \phi \). The identity eq. (B.9) can also be deduced from eq. (B.8) because

\[ \mathbb{E}_p \left[ \frac{1}{|\rho|} \int_{\square_m} |\nabla \phi^*(\mu, \cdot, m, p) - \nabla \phi^*(\mu, \cdot)|^2 \, d\mu \right] \leq 2 \mathbb{E}_p \left[ \frac{1}{|\rho|} \int_{\square_m} |\nabla \phi^*(\mu, \cdot, m, p) - \nabla \phi^*(\mu, \cdot)|^2 \, d\mu \right] \]

By eq. (4.26) and eq. (4.27), the second term can be bounded by \( J(\square_m, p, \mathbf{a}(\square_m)p) \); and by eq. (4.19) and eq. (B.2), this quantity converges to 0 as \( m \to \infty \). From now on, we thus focus on the proof of eq. (B.8).

The idea of the proof of eq. (B.8) is very close to that for eq. (B.7) and eq. (B.19). We fix a mescroscopic scale \( n = \left\lfloor \frac{m}{3} \right\rfloor \), and then use the decomposition

\[ \mathbb{E}_p \left[ \frac{1}{|\rho|} \int_{\square_m} |\nabla \phi^*(\mu, \cdot) - \nabla \phi^*(\mu, \cdot, m, p)|^2 \, d\mu \right] \leq 2 \mathbb{E}_p \left[ \frac{1}{|\rho|} \int_{\square_m} |\nabla \phi^*(\mu, \cdot)|^2 \, d\mu \right] \]

For the first term eq. (B.21)-a, we use the stationarity to transform to the integration on unit cube \( \square_0 \), and then use the fact that \( \nabla \tilde{\phi}_{p,n} \) converges to \( \nabla \phi \) in \( L^2_{\mu,\text{loc}} \) to get that

\[ \lim_{m \to \infty} \mathbb{E}_p \left[ \frac{1}{|\rho|} \int_{\square_0} |\nabla \tilde{\phi}_{p,n}(\mu, \cdot) - \nabla \phi(\mu, \cdot)|^2 \, d\mu \right] = 0. \]
Thus it suffices to finish the second term eq. (B.21)-b. We use the definition in eq. (B.6) and Jensen’s inequality to get that

eq (B.21)-b
\[ 2\mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{\Box_m} \left( \int_{y + \Box_n} \nabla \phi(\mu, y, x + \Box_n, p) \, dx \right)^2 \, d\mu(y) \right] \]
\[ \leq 2\mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{\Box_m} \int_{y + \Box_n} |\nabla \phi(\mu, y, x + \Box_n, p) - \nabla \phi(\mu, y, \Box_m, p)|^2 \, dx \, d\mu(y) \right] \]
\[ \leq 2 \times 3^{d-n} \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m+3^n} \setminus Q_{3^m-10^n}} \int_{y + \Box_n} |\nabla \phi(\mu, y, x + \Box_n, p) - \nabla \phi(\mu, y, \Box_m, p)|^2 \, dx \, d\mu(y) \right] \]
\[ \leq 2 \times \text{eq. (B.21)-b1} + \text{eq. (B.21)-b2} + \text{eq. (B.21)-b3} , \]
where in the last line we decompose once again the integration into three terms with respect to the domain

\text{eq. (B.21)-b1} = 3^{d-n} \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m-20^n}} \int_{y + \Box_n} |\nabla \phi(\mu, y, x + \Box_n, p) - \nabla \phi(\mu, y, \Box_m, p)|^2 \, dx \, d\mu(y) \right] ,
\text{eq. (B.21)-b2} = 3^{d-n} \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m+3^n} \setminus Q_{3^m-10^n}} \int_{y + \Box_n} |p + \nabla \phi(\mu, y, x + \Box_n, p)|^2 \, dx \, d\mu(y) \right] ,
\text{eq. (B.21)-b3} = 3^{d-n} \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m+3^n} \setminus Q_{3^m-10^n}} \int_{y + \Box_n} |p + \nabla \phi(\mu, y, \Box_m, p)|^2 \, dx \, d\mu(y) \right].

The terms eq. (B.21)-b2 and eq. (B.21)-b3 are easy to treat as they are boundary layer terms. For eq. (B.21)-b2 we can use the energy bound

\text{eq. (B.21)-b2} \leq C 3^{-(m-n)} \mu(\Box_n, p) \xrightarrow{m \to \infty} 0.

For eq. (B.21)-b3, since the function $|p + \nabla \phi(\mu, y, \Box_m, p)|^2$ does not involve $x$, we use Fubini’s lemma that

\text{eq. (B.21)-b3}
\[ = 3^{d-n} \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m-20^n}} \left( \int_{Q_{3^m+3^n} \setminus Q_{3^m-10^n}} 1_{\{x + \Box_n\}} \, dx \right) |p + \nabla \phi(\mu, y, \Box_m, p)|^2 \, d\mu(y) \right] \]
\[ \leq \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m+3^n} \setminus Q_{3^m-10^n}} |p + \nabla \phi(\mu, y, \Box_m, p)|^2 \, d\mu(y) \right] \]
\[ = \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m+3^n} \setminus Q_{3^m-20^n}} |p + \nabla \phi(\mu, y, \Box_m, p)|^2 \, d\mu(y) \right] + \mathbb{E}_\rho \left[ \frac{1}{\rho |\Box_m|} \int_{Q_{3^m+3^n} \setminus Q_{3^m}} |p|^2 \, d\mu(y) \right] \]
\[ \xrightarrow{m \to \infty} 0. \]

Here from the second line to the third line, we use the fact that the gradient contributes only on $Q_{3^m+3^n} \setminus Q_{3^m-20^n}$. Then we do a decomposition: the integration on $Q_{3^m+3^n} \setminus Q_{3^m}$ can be calculated directly, since $\phi(\mu, \Box_m, p)$ is $\mathcal{F}_{\Box_m}$-measurable and the gradient vanishes; the integration on $Q_{3^m+3^n} \setminus Q_{3^m-20^n}$ can be bounded by the boundary layer estimate in Lemma B.6.
Finally, we focus on eq. (B.21)-b1. We rewrite the integration $\int_{Q^{m-10\times 3^n}}$ as

eq (B.21)-b1

$$
\nu\left(\square_n, p\right) - \nu\left(\square_m, p\right) + 5 \times 3^{-2m} \Lambda |p|^2,
$$

where we used the quadratic response (4.9) from the second line to the third line. This implies that \( \lim_{m \to \infty} \) eq. (B.21)-b1 = 0, and thus completes the proof of eq. (B.8). □

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Figure 6. The function $\ell_{p,\mathbb{D}_m} + w_x$ is a sub-minimizer for the problem $\nu(\mathbb{D}_m, p)$, which combines the minimizer in cubes of scale $3^n$ biased by a vector $x$ (the cubes in blue), and a minimizer in $U_x$ (the domain in red).

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