INTRODUCTION

Cold field electron emission (CFE) is an electron emission regime where electrons are field emitted from a surface, by tunneling from states near the emitter Fermi level (Forbes, 2012; Forbes et al., 2015). The emission process is controlled by the height and width of the surface potential-energy (PE) barrier, and the barrier width is determined by the magnitude $F$ of the local barrier electrostatic field. For metal emitters, the emission current $I$ from a single tip emitter is given by a Fowler-Nordheim-type (FN-type) equation (Forbes, 2012; Forbes et al., 2015):

$$I = A_F a \phi^{-1} F^2 \exp\left[-v_F b \phi^{3/2} F\right],$$

where $a$ and $b$ are the first and second FN constants, $\phi$ is the local work-function, $v_F$ is the barrier-form correction factor for the particular potential-barrier shape assumed, and $A_F$ is the formal emission area, as defined in (Forbes et al., 2015). It is often assumed that this equation is also an adequate approximation for emission from non-metallic emitters.

Carbon nanotubes (CNTs) (Iijima, 1991) are comprised of graphene sheets rolled into seamless hollow cylinders with diameters ranging from 1 nm to about 50 nm (Saito et al., 1998). CNTs have the following properties that make them appropriate as CFE sources: (1) high aspect ratio, (2) small apex radius of curvature, (3) high chemical stability, and (4) high mechanical strength (Bonard et al., 1999; Choi et al., 1999; Gröning et al., 2000; Saito et al., 1998; Saito & Uemura, 2000). This study uses Nanocyl™ NC7000 (Nanocyl S. A., Belgium) thin multi-walled carbon nanotubes (MWCNTs) with 9.5 nm diameter and high aspect ratio (>150), fabricated by chemical vapor deposition (Bonard et al., 1998, 2002; de Jonge & Bonard, 2004).

Individual MWCNTs (Kim, 2013) were melt-welded into a fine glass capillary tube, using a "glass puller" of the form...
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shown in Fig. 1. This creates a prepared glass micro-point emitter (Mousa & Hibbert, 1993a) (Fig. 1), which will field emit when a sufficiently high electric field is applied. In this work a new kind of emitter is introduced which is based on internally a MWCNTs enclosed within a glass microemitter. The composite emitter is fabricated using a pulling technique. The analyses were essentially carried out using a field emission microscope. This system could be evacuated to a base pressure of about 10⁻⁹ mbar after baking at ~200°C overnight. The cathode was mounted ~10 mm away from the phosphorus screen and a current limiting resistor of 100 MΩ was used.

MATERIALS AND METHODS

As indicated, composite glass-on-MWCNTs emitters were prepared by the technique of pulling heated glass capillary tubes into fine points (Mousa & Hibbert, 1993a), using the device shown in Fig. 1. In this technique, two bearings are located accurately on plates supported by three stainless steel rods fixed rigidly to the frame of the control unit. This frame is strong enough to serve as a stable base for the instrument. A glass tube (outside diameter=1 mm, internal diameter=0.1 mm) fits inside these bearings, between the upper and lower chucks, with a furnace loop located around it. The MWCNTs were entered into the opposite end of each glass tube so they would protrude at the tip with a wire plunger (Mousa & Hibbert, 1993b). The lower chuck spindle can slide vertically, in order to pull down the glass tube under gravity. In operation, the temperature of the furnace is raised to the glass softening point (1400 K for borosilicate glass) (Mousa, 1991). Prepared emitters were investigated in a locally built point-to-plane field electron microscope (FEM), with an axial tip-to-screen distance of 10 mm (Al-Qudah et al., 2015; Mousa et al., 2015). The FEM was evacuated to ultra high vacuum conditions by the combination of a rotary pump (which produces a backing pressure of about 10⁻⁹ mbar) and an oil
diffusion pump system, connected to a liquid nitrogen (LN2) trap. Without baking, the FEM reaches a base-pressure of about 10^{-7} mbar. The pressure reaches about 10^{-9} mbar after baking the system for 12 hours at 180°C, and adding liquid nitrogen to the trap (Latham & Salim, 1987; Madanat et al., 2015; Moran Meza et al, 2015; Mousa et al., 2015). In operation, the FEM applied voltage \( V \) is increased slowly until a “switch-on voltage” \( V_{SW} \) is reached, at which point the emission current suddenly “switches on”. At \( V_{SW} \) the current increases rapidly from about a nanoampere to a much greater saturated value \( I_{SAT} \) as found previously (Latham & Mousa, 1986). The current-voltage (\( I-V \)) data yields a FN plot that has linear segments.

**RESULTS**

To obtain quantitative information about the microscopic properties of the emitting regime, the total electron emission current \( I \), has been plotted as a function of the applied voltage \( V \) between the emitter and the anode. Emission characterization will be presented as \( I-V \) plots and as the related FN plots. Throughout the experiments, electron emission images have been recorded by means of a digital camera, in order to study the spatial distribution and stability of the emission current.

Fig. 2A shows an optical image of the first tip to be discussed. Fig. 2B and C show the \( I-V \) characteristic and a related FEM image. As we increased the applied voltage slowly, a point was reached, at voltage \( V_{SW}=4,300 \) V, where the emission current suddenly 'switched-on' from an effective zero-value to a stable saturated value \( I_{SAT} \). Fig. 2B shows the switch-on at \( V_{SW}=4,300 \) V, \( I_{SW}=25 \) \( \mu \)A. As voltage is then reduced, the “saturated” region extends down to the values \( V_{SAT}=1,900 \) V, \( I_{SAT}=1 \) \( \mu \)A. Further voltage reduction causes the emission current to fall to nearly zero, reaching the threshold values \( V_{TH}=700 \) V, \( I_{TH}=6.0 \times 10^{-12} \) A. As shown in Fig. 3, once the emitter leaves saturation, the (decreasing) current-voltage relationship generates an approximately linear FN plot, with slope about \(-5,970 \) decade V.

Fig. 3. Fowler-Nordheim (FN) plot of tip CNT-A4 showing linear trend at low field region.

Fig. 4. (A) The \( I-V \) plot of this tip shows switch-on phenomena (\( V_{SW}=4,500 \) V, \( I_{SW}=23 \) \( \mu \)A). (B) Fowler-Nordheim (FN) plot. (C) Field electron microscope images after switch-on.
The experiment was then repeated after allowing “relaxation”. (Relaxation is the name we give to the procedure whereby, after an initial experimental run, we turn off the voltage, and leave the sample under vacuum for 12 hours before repeating an experimental run [Mousa et al., 2015]). During this time the pressure increased to $4.9 \times 10^{-8}$ mbar. The current-voltage characteristics show initiation by low current emission at $V=4,400$ V, $I=1.3$ nA. This is followed by switch-on at $V_{SW}=4,500$ V, $I_{SW}=23$ μA to a saturated current $I_{SAT}$ that may be seen in the expanded voltage range in Fig. 4. The emission characteristics of this tip show a saturation effect for high emission current. The dashed line represents the data for the low field region, which yields a linear FN plot with slope=$-10,470$ decade V as shows in Fig. 4C.

As presented in Fig. 4C, the emission image showed large bright spots, spread over most of the screen. We then repeated the experiment, for the same tip (CNT-A4) and under similar conditions (but with a baking period of 14 hours).

The following figures (Fig. 5-9) show cycles of increasing and decreasing the applied voltage.

The emission images shown in Fig. 6 (when increasing the applied voltage) are more concentrated. However, the image became large and diffuse when decreasing the voltage as shown in Fig. 7C.

In Fig. 9, the emission images show a stable bright spot, but when the applied voltage was decreased and reached the point $V_{SW2}=2,800$ V, $I_{SW2}=5.8$ μA the emission current increased suddenly. In Fig. 8B, the I-V plot shows evident suddenly increasing of emission current. Decreasing the applied voltage from 3,250 V to 1,000 V gave currents in the range 5.8 μA down to 9.9 pA.

The shift of the emission current to lower values may occur because of surface impurities; alternatively, this shift could result from the presence of multiple emitting sites on the tip that switch on at different voltages (Mousa & Hibbert, 1993b).

Studies were also carried out on a second tip (CNT-A5). The...
Fig. 7. Decreasing voltage. (A) $I$-$V$ plot (3,500–600) V. (B) Fowler-Nordheim (FN) plot showing saturation for high emission current, linear FN plot in the low field region, with slope equal to about $-4.010$ decade V. (C) Emission image ($V=3,200$ V, $I=13$ μA).

Fig. 8. Re-applied voltage again. (A) Emission image before switch-on ($V=3,050$ V, $I=0.6$ μA), decreasing voltage. (B) $I$-$V$ plot (3,250–1,000) V. (C) Fowler-Nordheim plot with slope about $-11,840$ decade V.
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optical micrograph in Fig. 10A shows the shape of this tip, which is a very smooth needle-like emitter. This tip has low switch-on voltage, as shown in its I-V characteristics and FN plot. With this tip, the voltage was cycled two times. As the applied voltage is slowly increased the emission current “switches on” from \( I = 6.0 \, \text{nA} \) to a stable saturated value of 3.2 \( \mu \text{A} \) to 1.0 \( \mu \text{A} \) by slowly reducing the voltage, the current decreases smoothly until it vanishes at a certain threshold voltage. Fig. 10B shows that for this emitter the switch-on occurs at \( V_{SW} = 2,300 \, \text{V} \), \( I_{SW} = 2.3 \, \mu \text{A} \). The saturated region extends down to \( V_{SAT} = 1,750 \, \text{V} \), \( I_{SAT} = 1 \, \mu \text{A} \). Beyond this, the emission current falls to nearly zero, as the applied voltage is decreased to a threshold value \( V_{TH} = 1,430 \, \text{V} \), with \( I_{TH} = 5.7 \, \text{pA} \). In Fig. 10C shows FN plot showing saturation for high emission current, linear FN plot at the low field region with slope about –8,720 decade V.

The emission images shown in Fig. 11 (when decreasing the applied voltage after switch-on) are more concentrated. The emission images in this case are small, concentrated bright and, stable. As shown in the related figures, a non-orthodox type of current-voltage behavior was observed with this emitter (R. Forbes; personal communication).

**DISCUSSION AND CONCLUSIONS**

It is helpful to attempt to understand the conditions under which continuous stable current may flow from the composite emitters. The electron emission dynamics is mostly controlled by the metal/glass interface (Mousa & Hibbert, 1993b). Our measurements show that MWNTs have high current emission. However, our observations suggest that the presence of localized energy states at the tip, instead of the metallic density of states encountered on usual emitters, influences greatly the emission characteristics. The deviations from FN behavior and saturation effects, as well as the presence of discrete current levels, observed on single MWNTs are characteristic of materials with non-metallic emitting states, as for example nanotips (Thien et al., 1996). Because of the nonmetallic nature of the tip, field penetration and fluctuations of the applied potential shift the position of the emitting levels. This could lead to the observed saturation at high applied voltages and strong instabilities of
the emission in some voltages ranges. The difference in energy between emitted and replacement electrons provokes a tip heating or cooling (Nottingham effect), related to the position of the Fermi level (Swanson & Bell, 1973; Thien et al., 1996). On tubes with emitting states just above the Fermi level this reduces the increase of the temperature with the current. As for the current steps, they most likely reflect changes in the surface structure (Swanson & Bell, 1973) and/or in the adsorbate coverage (adsorption, desorption or diffusion).
(Lee et al., 1995). The former possibility is unlikely because of the structure of the cap, whereas it is well known that small changes in adsorbate coverage result in large variations of the work function (Kyritsakis & Xanthakis, 2015; Lee et al., 1995). This effect is probably enhanced by the non-metallic character of the tip. Finally, since each tube has a unique tip, and hence a unique density of states at the tip, some tubes should be far more efficient emitters than others, depending on the position of the localized states with respect to the Fermi level. Some points are however not yet clarified, especially the behavior beyond the saturation point and the abrupt increase in the I-V characteristics some single MWNTs. The nature of the current saturation phenomenon with stable charge distribution on the emitter surface is not fully understood. Amongst other possible explanations, it could indicate a new field emission mechanism. Further research will be needed to clarify this point. Nonetheless, it appears that such tips might provide an electron source with stable supply mechanism of saturated current for technological application. All of the tips were able to serve rather high emission currents of up to 20 μA or more without being destroyed. This indicates protective properties of the glass microemitters with internal MWCNTs.

CONFLICT OF INTEREST

No potential conflict of interest relevant to this article was reported.

ACKNOWLEDGMENTS

Emad Bani Ali would like to thank the Scientific Research Support Fund (SRSF) for their support. The authors would like to thank Deanship of the Academic Research at Mu’tah University for supporting this work. Authors acknowledge Prof. Dr. Saadi Abdul Jawad for providing the Nanocly™ NC 7000 (MWCNTs).

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