Multiwavelength Observations of an Eruptive Flare: Evidence for Blast Waves and Break-out

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Abstract Images of an east-limb flare on 3 November 2010 taken in the 131 Å channel of the Atmospheric Imaging Assembly onboard the Solar Dynamics Observatory provide a convincing example of a long current sheet below an erupting plasmoid, as predicted by the standard magnetic reconnection model of eruptive flares. However, the 171 Å and 193 Å channel images hint at an alternative scenario. These images reveal that large-scale waves with velocity greater than 1000 km s\textsuperscript{-1} propagated alongside and ahead of the erupting plasmoid. Just south of the plasmoid, the waves coincided with type-II radio emission, and to the north, where the waves propagated along plume-like structures, there was increased decimetric emission. Initially the cavity around the hot plasmoid expanded. Later, when the erupting plasmoid reached the height of an overlying arcade system, the plasmoid structure changed, and the lower parts of the cavity collapsed inwards. Hot loops appeared alongside and below the erupting plasmoid. We consider a scenario in which the fast waves and the type-II emission were a consequence of a flare blast wave, and the cavity collapse and the hot loops resulted from the break-out of the flux rope through an overlying coronal arcade.

Keywords: Solar flare — coronal loops, magnetic field, flux rope, magnetic reconnection.

1. Introduction

The flare and flux rope eruption observed on the east limb of the Sun on 3 November 2010 showed many of the expected features of eruptive flare models (Carmichael, 1964; Sturrock, 1966; Hirayama, 1985; Kopp and Pneuman, 1976). In particular, a fast moving plasmoid, interpreted as a flux rope, appeared above the flare site with what looked like a long, hot current sheet below (Reeves and Golub, 2011; Cheng et al., 2011; Savage et al., 2012; Hannah and Kontar, 2012).

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Hot, inflowing plasma was observed below the plasmoid seemingly causing it to accelerate and drive a shock, indicated by type-II radio emission, in the low corona (Bain et al., 2012; Zimovets et al., 2012). The eruption also triggered transverse oscillations in the surrounding loop systems (White and Verwichte, 2012; White, Verwichte, and Foullon, 2012).

To understand the connection between the loops and the erupting plasmoid, we took a careful look at the relation between the hot (131 Å) and cooler (193 and 171 Å) plasma emission. This revealed fast waves along plumes outside the hot cavity that were also headed by enhanced radio emission, suggesting that the flare caused shock-like signatures considerably beyond the plasmoid, possibly due to a flare blast wave. Although blast waves at the onset of energetic flares are predicted (Vršnak and Lulić, 2000a, 2000b; Vršnak and Cliver, 2008; Magdalenić et al., 2008, 2010), they have proved difficult to observe. One major problem was that they propagate with speeds greater than 1000 km s\(^{-1}\), and therefore required high image cadence and a large field of view to be detected in images. Thus it was not until images from the Atmospheric Imaging Assembly (AIA; Lemen et al., 2012) became available that there was a reasonable chance of seeing them. Blast waves have however been invoked to explain the high (greater than 500 km s\(^{-1}\)) Doppler shifts observed in Fe \text{XIX} and Fe XX ultraviolet emission lines seen across a significant fraction of flaring active region coronae (Innes et al., 2001; Tothova, Innes, and Stenborg, 2011), and also in events exhibiting type-II radio bursts where there was only slow or no associated coronal mass ejection (CME) (Vršnak et al., 2006; Magdalenić et al., 2012).

The 3 November 2010 eruption is also significant because of its well-observed inflows below the erupting plasmoid which have been interpreted as inflows to a current sheet. As pointed out by Hannah and Kontar (2012), these inflows are hot (8–14 MK), not cool, as predicted by the standard model. If the fast waves result from a flare blast wave, then the plasmoid acceleration may have occurred before the inflows, and this leads us to consider the possibility that the inflows are a consequence of flux rope break-out through the overlying coronal arcade (Antiochos, DeVore, and Klimchuk, 1999; Karpen, Antiochos, and DeVore, 2012). Break-out leads to reconnection in the corona, the outward acceleration of the flux rope, and the formation of postflare loops on the side and below the break-out site, as well as inflows to a current sheet below the flux rope. We therefore take a further look at the structure of the flare heated plasma seen in the 131 and 94 Å images to understand the dynamics beyond the regions of the previously identified inflows and current sheet.

In this paper we first show the evidence for fast propagating waves and their relation to enhanced radio emission. We then compare the timing of break-out with the onset of plasma inflow below the plasmoid, and the structures seen in the 131 and 94 Å channel emissions. A summary sketch of the observed structures and our interpretation is given in the final section.

2. Observations and Results

The eruption occurred in active region NOAA 11121 which appeared on the eastern limb (N38E18) on 3 November 2010. The C4.9 flare was first seen in
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STEREO-B Extreme UltraViolet Imager (EUVI; Howard et al., 2008) images at 12:10 UT and had a hard X-ray onset at 12:13 UT and peak at 12:14 UT (Zimovets et al., 2012). The event, observed by the AIA onboard the Solar Dynamics Observatory (SDO; Pesnell, Thompson, and Chamberlin, 2012), the Reuven Ramaty High Energy Solar Spectroscopic Imager (RHESSI; Lin et al., 2002), the Nançay Radioheliograph (NRH; Kerdraon and Delouis, 1997), and the San Vito Radio Solar Telescope has caught attention because several of its features resemble the standard flare model and has presented a plausible case for a piston-driven shock (Bain et al., 2012; Zimovets et al., 2012).

In this analysis, we discuss AIA 131, 94, 193, and 171 Å image sequences taken with a cadence of 12 s. Most of the images are shown after the subtraction of a pre-flare image taken at approximately 12:10 UT (i.e. base-difference images) to reveal the changes caused by the flare. To find the location of the radio emission, we have used high-cadence radio images of the Sun from the NRH, which operates in the frequency range of 160 – 435 MHz and probes the coronal plasma at heights $\leq 1.5 R_\odot$ (Kerdraon and Delouis, 1997).

An overview of several of the features discussed in previous papers is shown in Figure 1. The 131 Å image, dominated by Fe XXI emission, shows a hot, $\approx 10$ MK, plasmoid directly above and connected to the flare site by a long, thin thread, which has been interpreted as a current sheet (CS). The plasmoid swept up the surrounding plasma into a bright rim of cooler, $\approx 2$ MK, 193 Å emitting plasma (Reeves and Golub, 2011; Cheng et al., 2011). Radio observations revealed a type-II radio burst near the leading edge of the expanding plasmoid system (Bain et al., 2012; Zimovets et al., 2012). Behind the plasmoid, converging towards CS, were hot plasma inflows (Savage et al., 2012). The ripples along the northern flank of the hot cavity have been interpreted as Kelvin-Helmholtz (KH) roll-ups (Foullon et al., 2011). Important features for our later discussion are also illustrated in the 193 Å (Figure 1b) and 171 Å (Figure 1c) base-difference images. Here dark represents pre-flare active region structures that have disappeared due to heating or eruption. For example, the dark loop in the 193 Å image corresponds to a pre-flare coronal loop or arcade. The 171 Å base-difference image shows that there was enhanced emission both south and north of the hot, plasmoid cavity.

2.1. Early EUV Waves

During the early stage of the flare/flux-rope eruption, fast waves were observed in sequences of 171 Å images. Figure 2 is a composite of AIA 171 Å and 131 Å base-difference images showing the regions and directions of the fast waves. This is one frame from the movie linked to the figure in the online version. In Figure 2, the EUV wave, south to the plasmoid is visible as enhanced 171 Å emission extending beyond the height of the 131 Å plasmoid. The enhanced emission seen in the north at this time is probably due to resonance scattering of extreme ultraviolet (EUV) flare emission because the whole region brightened simultaneously and propagating waves were not seen here until 12:15:00 UT. The wave speeds can be deduced from the space-time images shown in Figure 3. This method computes the time when there is an abrupt change in the intensity at each spatial position.
Then a linear function is fitted to the values to obtain the velocity. When the front changes speed, it is sensitive to the start and end distance used to fit the wave front. For example, the front in Figure 3b changed from 1200 km s\(^{-1}\) at the start to 700 km s\(^{-1}\) at the end. On the image, we draw the line corresponding to the best estimate of the velocity and provide the velocity range. The waves to the south had a very high plane-of-sky speed, about 2200 km s\(^{-1}\). The speed was about twice that of the plasmoid and the waves along the plumes to the north. Both the southern waves and the plasmoid eruption started at 12:13 UT, a couple of minutes before the waves to the north. The composite AIA 131 and 171 Å base difference movie clearly shows the bright EUV wavefront south of the 131 Å plasmoid, and the propagation of the EUV wave along the plumes, north of the flare site.

### 2.2. Radio Emission

The radio emission has been discussed in the context of the RHESSI hard X-ray emission by Zimovets et al. (2012) and Bain et al. (2012). At flare onset there was a short burst of hard X-ray (25–50 keV) and microwave (3000–5000 MHz) emission, peaking at \(\approx 12:14\) UT. A few minutes later, there was a decimetric type-II radio burst that showed distinct band-splitting, indicating two outward-moving sources. Zimovets et al. (2012) interpreted these as the up- and downstream regions of a shock but it is also possible that the splitting is due to waves along two different density structures. Zimovets et al. (2012) found that one source started slightly off-set (100 Mm) south of the hot plasmoid trajectory, and had a speed 2240 km s\(^{-1}\), whereas the other was in line with the plasmoid and had a velocity 1500 km s\(^{-1}\).

In Figure 4, we compare the timing and sites of the radio emission with the waves. Figure 4a, shows that early on there was enhanced 445 and 432 MHz emission across a large portion of the active region, with the most intense emission on either side of the erupting plasmoid. The southern radio source...
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Figure 2: Composite of 131 Å (cyan) and 171 Å (red) base difference images showing a broad enhanced front of 171 Å emission just south of the hot plasmoid, observed at 131 Å. The white lines show the position of the space-time images in Figure 3. This is one frame from the movie 131_171.avi.

is the low frequency component tracked by Zimovets et al. (2012) and estimated to have a velocity 2240 km s\(^{-1}\). As can be seen in the images, it coincided with the trajectory of the fast, 2200 km s\(^{-1}\) waves inferred from the 171 Å images, so it is possible that the waves excited this component of the radio emission.

The second component which was labelled the high frequency component by Zimovets et al. (2012), was seen along the trajectory of the hot plasmoid. It did not appear until 1 min later at 327 MHz (Figure 4b), and had a speed of 1500 km s\(^{-1}\), which was faster than the hot plasmoid plane-of-sky speed, and close to the speed of the 211 Å leading edge (Zimovets et al., 2012). Subsequently, emission in front of the plasmoid dominated at frequencies at and below 327 MHz, and both high and low frequency radio components appeared over an extended region above the plasmoid. The radio emission may then have been produced by up- and downstream regions of a shock or by waves along different structures as suggested by earlier images.

In Figure 4, we also show 150 MHz contours. This emission comes from waves in lower density plasma than the higher frequency emission, so it is seen later and higher in the corona when caused by an outward propagating source. On this day, there was a 150 MHz frequency source at the head of the northern plumes (Figure 4a) where the northern 171 Å fast waves were observed. This source had been there long before the flare and it brightened at the time of the flare (Figure 4b). The NRH maps at the time of the early 150 MHz source brightening are shown in Figure 4b. The center of the 150 MHz emission moved...
Figure 3. Space-time base difference images of 171 Å (top row), 193 Å (middle) and 131 Å (bottom) emission, along the three white lines shown in Figure 2. The left column is along the northern line, the middle along the plasmoid and the right column along the southern line. Dashed white lines show the track for a front with the labeled speed.

in front of the hot plasmoid at 12:16:24 UT, as part of the type-II burst. Later, as the flare evolved the centroid of the low frequency (150–327 MHz) emission moved south (Figures 4d and 4e). The location of the radio sources with respect to the eruption site is shown in Figure 4e, where we have overlaid radio source (150, 228, and 327 MHz) contours on the AIA 193 Å running difference image of the eruption at 12:15:19 UT. This image shows that from 12:14-12:19 UT, the radio sources drifted southward (indicated by an arrow). Figure 4f displays the
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2.3. Break-out

Break-out occurred when the flux-rope broke through the overlying cavity system. Several structural changes appeared at the time of break-out. The overlying loop system disrupted and the plasmoid structure changed. At the same time, a hot thread of emission appeared below and attached to the plasmoid, and the earlier expansion of the cavity changed to contraction.

Some of these features are illustrated in Figure 5. These are time-distance images taken across (top) and along (bottom) the direction of propagation of the erupting flux rope, as shown in Figure 6b. These two perpendicular directions have been chosen because they allow comparison of the expanding and contracting cavity below the plasmoid and the outward propagating waves. The first dashed line marks the flare onset time, and the second dashed line the time of break-out, as deduced from the change in structure of the plasmoid leading edge, shown in the bottom row of Figure 5.

The images at 171 and 193 Å show the broad bright front of the EUV waves generated at onset, followed by dimming north and south of the flux rope (Figure 5a). Loop brightening and oscillations appeared immediately behind the front over most the central region. The early wave fronts are not visible in images of the hotter, 131 Å emission. The 131 Å image on the top row shows the expansion and contraction of the flux rope cavity, especially in the north, where the cavity edge lies along the inner edge of the 193 Å faint front indicated with an arrow in Figure 5b. At the time of the second dashed line the cavity expansion changes to contraction.

As mentioned above, the second dashed line in Figure 5 was placed to coincide with the time that the plasmoid changed shape. The leading front in the 193 and 131 Å space-time images on the bottom row split at the time of the second dashed line: the plasmoid slowed down while the fast front seen in 193 Å images continued through the corona. At this time, the 193 Å loop opened up and the bright 131 Å flux rope changed from circular to concave (Cheng et al., 2011; Zimovets et al., 2012), and it looks as though there was a fundamental change in the plasmoid propagation. We associate this sudden change with break-out. We also note that according to Figure 5f, the current sheet formed about 30 s (two frames) before the second dashed line. Thus the current sheet formed while the cavity was still expanding.

The structural changes are best seen in the accompanying movies which are two colour composites of the 131 and 94 Å, and the 131 and 171 Å emissions. A few frames from the movies are shown in Figure 6. The top two rows show the emission at onset and the last row shows the hot post-eruption loops as seen in the 131/94 Å images. The selected images show that initially (around 12:13:24 UT), a bright tongue of 94 Å emission appeared above the 131 Å flare emission and 171 Å waves. The tongue then separated and was filled with the plasmoid or flux rope of hotter 131 Å emission. In the image shown in Figure 6d,
Figure 4. Evolution of radio emission in the context of the EUV waves and hot plasmoid. The bottom panels show the drifting type II sources indicated by arrows.
the 131 Å plasmoid appears to be encased in a cocoon of 94 and 131 Å emission. On the northern edge of the cocoon, close to the limb bright 94 Å emission appeared. This seems to have been early, low-lying loops in the bright flare arcade seen in Figure 6. On the south of the plasmoid cavity, a region of faint, hot emission is just visible. The 131 Å image at 12:14:45 UT was also shown in White, Verwichte, and Foullon (2012), where they comment on the faint line of hot emission, indicated by an arrow, running from the plasmoid to the site.
Figure 6. Evolution of hot plasma in the context of oscillating loops seen in 171 Å channel. (a and b) 171 Å and 131 Å composites; (c - f) 94 Å and 131 Å composites. On panel (b) the white lines show the position of the time-distance images in Figure 5. The top row shows frames from the movie 131_171.avi and the other four frames are from the movie 131_94.avi.

of the hot loop seen later in 131 Å images, and shown in Figure 6. The 94 Å emission, especially the emission seen at 12:25:02 UT, suggests that this line of 131 Å emission outlined the top of a hot-loop arcade that formed at the time of break-out.

2.4. STEREO View of the Flare Site

This flare that was seen over the east limb from Earth, was observed near disk center in the STEREO-B/EUVI (hereafter EUVI-B) images. Savage et al. (2012) have identified the main flare footpoints and flaring loops in the EUVI-
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Figure 7. Co-temporal images of AIA 171 and STEREO-B 171 Å for identification of loops and plumes. Crosses of the same colour identify the same points in the two images.

B 195 Å images. At 195 Å, the cool plumes and loops are less visible than at 171 Å. Therefore, in Figure 7 we show the 171 Å EUVI-B and AIA images. Unfortunately, the EUVI-B image is partially saturated by the flare emission, so that loops that were below and along the line-of-site to the erupting flux rope cannot be seen. However the loops that were to the north and south, and the sites of the plumes are clearly visible.

To find the positions of structures from the different perspectives, we have used the SolarSoft routine scc_measure.pro (Thompson, 2006). We selected a point on the AIA image and then identified the same feature on the epipolar line drawn on the EUVI-B image. The same positions are marked with the same color cross in the two images. The EUVI image shows that the plumes were more than 100 arcsecs from the flare site. We note that there was an arcade of large loops to the south and a faint (in 171 Å) arcade of loops to the north linking the flare footpoints, as identified by Savage et al. (2012).

3. Summary and Discussion

We have taken a second look at the waves and structures observed by SDO/AIA during a flare eruption that occurred on 3 November 2010, on the east limb of the Sun as seen from Earth. At flare onset, fast 1000−2000 km s\(^{-1}\) EUV waves were seen propagating along cool (1 MK) plume-like structures that were situated about 100 arcsec from the flare site, and headed by decimetric emission. These waves were traveling ahead and alongside a fast, flare-heated plasmoid seen in 131 Å images. Loops ahead of the plasmoid, heated to about 5 MK, were seen in the 94 Å channel images. Figure 8, is a simplified sketch of the various features.

When the hot plasmoid reached the apex of the pre-flare loops, the leading edge of the plasmoid accelerated away from the main core. At the same time,
hot loops were seen below and on the side of the plasmoid and the region below collapsed inward to fill the space evacuated by the outward moving plasmoid, as sketched in Figure 8.

Initially flaring in the low corona/chromosphere caused a rapid pressure increase, flux rope destabilization, cavity expansion, and high energy particles. Waves propagated outward in all directions. Heating was confined to the inner flux-rope structure and its cavity. The rapidly rising flux rope soon encountered the overlying arcade. Reconnection with the overlying field led to break-out [Antiochos, DeVore, and Klimchuk, 1999; Karpen, Antiochos, and DeVore, 2012]. The formation of high-lying, hot loops alongside and just below the plasmoid and the plasma inflows were a consequence of break-out. Post-flare loops appeared below the plasmoid, close to the limb. Since the eruption occurred over the limb as seen from Earth, it is not possible to determine the exact flare site, the magnetic field configuration, or the flare trigger. However, the STEREO 171 Å image (right panel of Figure 7) shows the complex quadrapolar magnetic field configuration associated with a possible null point above the flare site, which are needed for the breakout reconnection.

This scenario is very different from previous ones for this event in two respects. First, we discovered fast waves at the onset of the flare and this led to the interpretation that a blast wave rather than a piston-driven shock caused the type-II radio emission. Second, we noted that the inflows occurred when the plasmoid structure changed and the over-lying 193 Å loop structure opened up and thus conclude that the inflows are as a consequence of break-out. In this scenario, the bright 131 Å thread that was below the plasmoid could be either the current sheet or a newly-formed loop seen edge on. The observation of hard X-ray emission near the plasmoid (Bain et al., 2012) favour the current sheet interpretation but the width of the emission and its stability suggest that it may be a newly formed loop similar to those seen on either side of the plasmoid.
We could not observe the formation of a piston-driven shock in the AIA 211 and 193 Å running difference images as observed in other flux-rope eruption events [Ma et al., 2011; Kozarev et al., 2011; Gopalswamy et al., 2012]. The flux-rope is heavily decelerated (mean speed ≈241 km s\(^{-1}\)) in the interplanetary (IP) medium, and there was no piston-driven shock in the IP space (e.g., IP type-II radio burst). Moreover, there is no strong lateral expansion of the flux rope, that could be sufficient to drive the shock at the flanks of the CME. [Patsourakos and Vourlidas, 2012]. We observed fast waves simultaneously at both sides of the flux-rope (along the plumes), therefore it is also possible that the pressure pulse overran the flux rope causing acceleration of the plasmoid in the low corona. Flare-heated plasma was observed at the flare site, in the flux-rope and along loops outlining the surrounding cavity, suggesting significant heating throughout the flux rope cavity, not just along the central current sheet (see movie 131_94.avi).

In conclusion, we reported a unique multiwavelength observation of the fast propagating EUV waves in AIA 171 Å and associated radio emissions ahead of these waves, which supports the flare blast-wave scenario as a possible candidate to generate high frequency type-II radio bursts. However, further analysis of such events will be helpful to understand the eruption processes associated with the generation of these fast waves.

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