Ambipolar Carrier Mobility in Binary Blend Thin Film of Non-Peripheral Alkylphthalocyanines

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Abstract. Charge carrier mobility in blend films of two types of soluble phthalocyanine derivatives with different substituent length, 1,4,8,11,15,18,22,25-octapentyl-phthalocyanine (C5PcH2) and 1,4,8,11,15,18,22,25-octahexyl-phthalocyanine (C6PcH2) has been investigated. The charge carrier mobility was measured by the time-of-flight technique. In the case of C5PcH2 blend ratio below 25 mol%, the high mobility, such as 0.8 - 1.1 cm²V⁻¹s⁻¹ for hole and 0.6 cm²V⁻¹s⁻¹ for electron, were obtained at room temperature. In the thin films with C5PcH2 above 30 mol%, the charge carrier mobility decreased by one order of magnitude and had the different temperature dependence from that below 25 mol%. The marked change of the charge carrier mobility depending on the blend ratio of phthalocyanine derivatives has been discussed by taking the miscibility and the molecular packing structure into consideration.

1. Introduction

Organic electronics is one of the most attractive and important research fields to realize light-weight, flexible and ubiquitous electronic devices in commercial base. In particular, a recent trend of this research field is “printed electronics” based on a solution processing, and realization of the low-cost organic thin film devices, such as organic light-emitting diodes, organic field effect transistors and organic photovoltaics, is promising. Required characteristics of organic semiconductors available for the printed electronics are high solubility in common organic solvents, spontaneous molecular alignment during the solvent evaporation to form a high quality thin film on a substrate, and high carrier mobility, and so on. Some organic semiconductor materials, such as pentacene derivatives [1,2], benzo[15]thienobenzothiophene derivatives [3] and dinaphthothienothiophene derivatives [4], were reported and must be the candidates applicable to the printed electronics.

Recently, we reported unique properties of a phthalocyanine derivative, 1,4,8,11,15,18,22,25-octahexylphthalocyanine (C6PcH2) [5]. C6PcH2 is soluble in common organic solvents and has the discotic liquid crystalline (LC) properties. C6PcH2 forms a hexagonal columnar structure by self-organization in the LC and crystal phases, and demonstrates the high ambipolar carrier mobility comparable to that of amorphous silicon, 1.08 and 0.43 cm²V⁻¹s⁻¹ for hole and electron at room temperature, respectively. C6PcH2 is available for bulk-heterojunction (BHJ) solar cells as a donor
material and the high power conversion efficiency (PCE) of 3.2 % [6,7] were reported. In the C6PcH2-based BHJ solar cells, the PCE could be improved to be 3.8% [8] using the blend donor materials of C6PcH2 and its analogue molecule, 1,4,8,11,15,18,22,25- octapentylphthalocyanine (C5PcH2).

Utilizing the mixture of small molecules with different alkyl-substituents, improvement of charge carrier mobility, and molecular alignment control have been reported [9-11]. The details of fundamental properties of the blend thin film of C5PcH2 and C6PcH2 should include an important essence to realize high performance devices, however, have not been clarified yet.

In this study, the charge carrier mobility in the blend films of C5PcH2 and C6PcH2 was investigated by a time-of-flight technique and discussed by taking the miscibility and molecular packing structure in the films.

![Figure 1. Molecular structures of C5PcH2 (n=5) and C6PcH2 (n=6).](image)

2. Experimental

The molecular structures of C5PcH2 and C6PcH2 are shown in Fig. 1. C5PcH2 and C6PcH2 were synthesized according to the literature [12] with slight modifications, and fully purified by column chromatography (silica-gel with toluene as eluent) followed by repetitive recrystallization from toluene-methanol (1:2) solution.

The mixtures of C5PcH2 and C6PcH2 were prepared by the following procedures. C5PcH2 and C6PcH2 were dissolved in toluene, and toluene was vaporized under a N2 atmosphere. The molar ratios of C5PcH2 and C6PcH2 were 12.5:87.5, 25:75, 30:70, 37.5:62.5, 50:50 and 75:25.

A sandwich cell, consisting of two indium tin oxide (ITO)-coated glass substrates and a 9-μm-thick PET spacer, was used for the TOF measurements. The actual cell gap was evaluated by an interference technique of light transmittance. The mixture of C5PcH2 and C6PcH2 was heated to a temperature in the isotropic phase for injecting into the cell by a capillary action, and cooled down to room temperature in a vacuum.

The phase transition behavior of the mixtures of C5PcH2 and C6PcH2 was investigated by differential scanning calorimetry (DSC, TA instruments Q2000) and polarizing optical microscopy (POM, Nikon ECLIPSE E600 POL) equipped with a Mettler FP82HT Hot Stage and Mettler FP90 Central Processor. X-ray diffraction (XRD) pattern was measured using an X-ray diffractometer (Rigaku RINT 2000) with Cu Kα line. A measured sample was prepared by separating the pair of glass substrates of the
sandwich cell. The lattice parameter \( d \) was calculated from the XRD patterns using Bragg’s law; \( n\lambda = 2d\sin\theta \), where \( n \) is an integer, \( \lambda \) is the X-ray wavelength of 1.5418 Å, and \( \theta \) is the Bragg angle.

3. Results and Discussion
In order to confirm the miscibility of C5PcH2 and C6PcH2, the thermal properties in the mixture of C5PcH2 and C6PcH2 were investigated. Figure 2 shows the phase transition diagram determined by using the DSC and POM of the mixture of C5PcH2 and C6PcH2. A LC phase appears between 161 °C and 169 °C for C6PcH2, while no LC phase exists in the case of C5PcH2. The clearing points of the mixtures were kept constant in the case of C5PcH2 blend ratio lower than 50 mol%. The melting points almost coincided with Le Chatelier-Schröder relation [13-15] exhibited as a theoretical thermal behavior of the melting point for each blend ratio. The eutectic point of the mixture could be approximately determined to be 20 mol% of the C5PcH2 blend ratio from the calculation using Le Chatelier-Schröder relation. Therefore, these results indicate that the mixtures of C5PcH2 and C6PcH2 possess the miscibility, and each compound is mixed in the molecular level.

![Figure 2. Phase transition diagram for the mixtures of C5PcH2:C6PcH2.](image)

Figure 3 shows typical decay curves of the transient photocurrents for hole and electron in the blend film of C5PcH2:C6PcH2=50:50 under an applied electric field of 27.9 kV/cm at room temperature. The decay curves with the clear kink points were observed throughout the temperature range of the crystal phase and the transit times for hole and electron transport could be determined. The decay curves of other blend ratios of C5PcH2 and C6PcH2 were also nondispersive behaviors except for those of the electron in the blend ratios of 30:70 and 75:25. The charge carrier mobility was independent of the applied field, which is a typical behavior observed in liquid crystalline semiconductors [16].

Figures 4(a) and 4(b) show the temperature dependence of the hole and electron mobility in the blend films of C5PcH2:C6PcH2, respectively. Comparing the temperature dependence, the obvious different behaviors could be found among the three regions, that is, in the C5PcH2 blend ratio below 25 mol%, 30-37.5 mol%, and 50-75 mol%. In the region of C5PcH2 blend ratio below 25 mol%, the charge carrier mobility was independent of the temperature and approximately 0.25 cm²V⁻¹s⁻¹ for hole and 0.3 cm²V⁻¹s⁻¹ above about 140 °C, while below about 140 °C the charge carrier mobility increased from 0.05 to 0.05 cm²V⁻¹s⁻¹. With decreasing the temperature. In the region of 50-75 mol%, the charge carrier mobility was independent of the temperature and approximately 0.25 cm²V⁻¹s⁻¹ for hole and 0.3 cm²V⁻¹s⁻¹ above about 140 °C, while below about 140 °C the charge carrier mobility increased from 0.05 to
0.1 cm²V⁻¹s⁻¹ with decreasing the temperature. The temperature dependence of the charge carrier mobility in crystal phase could be explained by the volume expansion of the unit lattice in the molecular packing structure. The detailed studies about the temperature dependence of the lattice parameter in the C5PcH₂:C6PcH₂ blend single crystal by using X-ray structure analysis are now in progress, and the relationship between the mobility and lattice parameter would be reported elsewhere.

Figure 3. Transient photocurrent decay curves of the blend film of C5PcH₂:C6PcH₂=50:50 under the field of 27.9 kV/cm at room temperature.

Figure 4. Temperature dependence of hole mobility (a) and electron mobility (b) in the C5PcH₂:C6PcH₂ blend films.

Comparing the charge carrier mobility in the blend films of C5PcH₂:C6PcH₂, the C5PcH₂ blend ratio dependence of the mobility at room temperature is shown in Fig. 5. In the case of C5PcH₂ blend ratio below 25 mol%, both hole and electron mobilities kept high values close to those of C6PcH₂. The hole mobility slightly decreased from 1.1 to 0.8 cm²V⁻¹s⁻¹ with increasing the blend ratio of C5PcH₂, while the electron mobility kept constant, 0.6 cm²V⁻¹s⁻¹. In the case of C5PcH₂ blend ratio above 30 mol%, both hole and electron mobilities markedly decreased by one order of magnitude, and the hole mobility slightly increased with increasing the blend ratio of C5PcH₂.

Figures 6(a) and 6(b) show the XRD patterns of the blend films with C5PcH₂:C6PcH₂ and the magnified figure in the range of 2θ=4-6°, respectively. In the case of C5PcH₂ blend ratio below 25
mol%, the diffraction peaks at 4.86° and 26.02° appeared, which correspond to the $d(100)$ lattice parameter of hexagonal column formed by C6PcH2 (Fig. 7(a)) and the $d(001)$ lattice parameter of inter-molecular distance, respectively. It is considered that the molecular packing structure like C6PcH2 was formed and the molecules were stacked in order, since C6PcH2 was dominant for the self-organization in the appearance of the liquid crystal phase in the C5PcH2 blend ratio below 25 mol%, therefore, the charge carrier mobility was kept high.

**Figure 5.** C5PcH2 blend ratio dependence of charge carrier mobility in the C5PcH2:C6PcH2 blend films at room temperature.

In the C5PcH2 blend ratio of 50 mol%, the two diffraction peaks at 4.74° and 4.89° were observed, which correspond to the $d(200)$ and $d(110)$ lattice parameter of rectangular column (Fig. 7(b)). It is supposed that the lattice constant in the $a$-axis direction of rectangular lattice became longer, while that in the $b$-axis direction became shorter, because equivalent mixing caused the distortion of the hexagonal lattice.

In the case of C5PcH2 blend ratio above 50 mol%, the diffraction peak shifted from 4.89° to 5.28° corresponding to the $d(100)$ in Fig. 7(a) as the blend ratio of C5PcH2 increased. It is considered that the mean length of the alkyl chains around the column became shorter with increasing the blend ratio of C5PcH2, because of short substituent of C5PcH2, resulting in the decrease of the inter-columnar distance.

**Figure 6.** XRD patterns of the C5PcH2:C6PcH2 blend films at room temperature in the angle ranges of 2-40° (a) and 4-6° (b).
The diffraction peaks around 26°, which originates from the molecular stacking, however, disappeared in the higher C5PcH2 blend ratio than 25 mol%, therefore, the marked decrease of the carrier mobility in the film with C5PcH2 above 30 mol% must be explained by the disordered stacking in the columnar structure. Since the inter-molecular overlapping of π-orbitals must strongly contribute to the carrier transport, therefore, it is considered that the carrier transport in the blend film of C5PcH2 and C6PcH2 should be dominated by the order of molecular packing like π-stacking rather than the lattice parameter like inter-columnar distance.

![Figure 7](image)

**Figure 7.** Schematic diagrams of (a) hexagonal and (b) rectangular columnar structures in the blend films of C5PcH2:C6PcH2.

### 4. Conclusions

The charge carrier mobility in the blend films of C5PcH2 and C6PcH2 were investigated. The phase diagram of the mixture of C5PcH2 and C6PcH2 was determined from the thermal properties, it was found that the mixture possess the miscibility. In the case of C5PcH2 blend ratio below 25 mol%, both hole and electron mobilities were kept high values comparable with those of C6PcH2, while both hole and electron mobilities markedly decreased by one order of magnitude in the case of C5PcH2 blend ratio above 30 mol%. From the results of XRD measurement, the molecular structure of the mixture film obviously changed depending on the blend ratio of C5PcH2, that must be strongly related with the behavior of the charge carrier transport.

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### References

[1] Chang J, Sakanoue T, Olivier Y, Uemura T, Dufourg-Madec M, Yeates S G, Cornil J, Takeya J, Troisi A and Sirringhaus H 2011 Phys. Rev. Lett. 107 066601

[2] Llorenete G R, Dufourg-Madec M, Crouch D J, Printchard R G, Ogier S and Yeates S G 2009 Chem. Commun. 3059

[3] Uemura T, Hirose Y, Uno M, Takimiya K and Takeya J 2009 Appl. Phys. Express 2 111501
[4] Nakayama K, Hirose Y, Soeda J, Yoshizumi M, Uemura T, Uno M, Li W, Kang M J, Yamagishi M, Okada Y, Miyazaki E, Nakazawa Y, Nakao A, Takimiya K and Takeya J 2011 Adv. Mater. 23 1626
[5] Miyake Y, Shiraiwa Y, Okada K, Monobe H, Hori T, Yamasaki N, Yoshida H, Cook M J, Fujii A, Ozaki M and Shimizu Y 2011 Appl. Phys. Express 4 021604
[6] Hori T, Miyake Y, Yamasaki N, Yoshida H, Fujii A, Shimizu Y and Ozaki M 2010 Appl. Phys. Express 3 101602
[7] Hori T, Fukuoka N, Masuda T, Miyake Y, Yoshida H, Fujii A, Shimizu Y and Ozaki M 2011 Sol. Energy Mater. Sol. Cells 95 3087
[8] Fukui H, Nakano S, Uno T, Dao Q, Saito T, Fujii A, Shimizu Y and Ozaki M 2014 Org. Electron. 15 1189
[9] Wegewijs B R and Siebbeles L D A 2002 Phys. Rev. B 65 245112
[10] Shimizu Y, Matsuda Y, Nekelson F, Miyake Y, Yoshida H, Fujii A and Ozaki M 2009 Proc. of SPIE 8279 82790G
[11] G. Schweicher, G. Gbabode, F. Quist, O. Debever, N. Dumont, S. Sergeyev and Y. H. Geerts: 2009 Chem. Mater. 21 5867
[12] Swarts J C, Langner E H G, Krokeide-Hove N and Cook M J 2001 J. Mater. Chem. 11 434
[13] Le Chatelier H 1894 C. R. Acad. Sci. 118 638
[14] Schröder I 1893 Z. Phys. Chem. 11 449
[15] Van Laar J 1903 Arch. Neerl. 8 264
[16] Yoshino K, Nakayama H, Ozaki M, Onoda M and Hamaguchi M 1997 Jpn. J. Appl. Phys. 36 5183