Strain Gated Bilayer Molybdenum Disulfide Field Effect Transistor with Edge Contacts

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Silicon nitride stress capping layer is an industry proven technique for increasing electron mobility and drive currents in n-channel silicon MOSFETs. Herein, the strain induced by silicon nitride is firstly characterized through the changes in photoluminescence and Raman spectra of a bare bilayer MoS2 (Molybdenum disulfide). To make an analogy of the strain-gated silicon MOSFET, strain is exerted to a bilayer MoS2 field effect transistor (FET) through deposition of a silicon nitride stress liner that warps both the gate and the source-drain area. Helium plasma etched MoS2 layers for edge contacts. Current on/off ratio and other performance metrics are measured and compared as the FETs evolve from back-gated, to top-gated and finally, to strain-gated configurations. While the indirect band gap of bilayer MoS2 at 0% strain is 1.25 eV, the band gap decreases as the tensile strain increases on an average of ~100 meV per 1% tensile strain, and the decrease in band gap is mainly due to lowering the conduction band at K point. Comparing top- and strain-gated structures, we find a 58% increase in electron mobility and 46% increase in on-current magnitude, signalling a benign effect of tensile strain on the carrier transport properties of MoS2.

Strain is a critical ingredient in modern transistor scaling. For Intel process technologies, the electron mobility decreased from 400 to 120 cm2/Vs when the industry migrated from 0.80 μm to 0.13 μm technology node due to the large vertical electric field1. Strain engineering has proven an effective route for mobility enhancement by modifying the carrier effective mass and mean scattering time2. For transistors with sub-100 nm gate length, strained silicon increased the saturated MOSFETs drive currents by 10–20% and mobility by >50%, and later ramped into high volume manufacturing on high performance microprocessors in the 90 nm logic technology3. The electron and hole mobility of silicon responds differently to externally applied stress. Longitudinal tensile stress along transistor channel improves electron mobility but degrades hole mobility4–6. p-MOSFET features a compressively strained SiGe film embedded in the source and drain regions7. The mismatch in the SiGe to Si lattice causes the channel to be under a uniaxial compressive strain, leading to significantly improved hole mobility8. For n-MOSFET, a post-salicide tensile silicon nitride capping layer was deposited on top of the transistor gate, wrapping both the gate and source drain area7. As a tensile stressor inside the nitride film tends to shrink, the stressor on the source and drain pulls apart the ends of the transistor channel and mainly produce a longitudinal tensile strain in the n-MOSFET channel2.

For three-dimensional semiconductor, the ultimate strain exertion is limited by both bulk defects and surface imperfections9. Transition metal dichalcogenides (TMDCs) are more suitable for strain engineering for two reasons: negligible bulk defects because of a thickness of only a few atomic planes and a fully-terminated surface, eliminating fab processing steps that usually targeted at the passivation of dangling bonds. A breaking strain up to 11% for MoS210 has been proved experimentally, whereas bulk silicon can be strained only 1.2% before fracture9. Theoretical studies have predicted that when an external tensile stress is applied, the electronic structure of monolayer MoS2 undergoes a series of variations: first, a direct-to-indirect band gap transition when the lattice constant is just slightly lengthened; second, a more drastic semiconductor-to-metal transition when the lattice constant is increased by more than 9.8%11. In particular, the tensile strain reduces the gap energy and the effective masses while the compressive strain enhances them11,12. Majority of experimental demonstration of the
strain effect on MoS$_2$ employs standard three-point or four-point bending apparatus together with micro-Raman facilities. Ultra-thin MoS$_2$ samples are firstly exfoliated and then clamped onto a bendable material such as polydimethylsiloxane (PDMS)\textsuperscript{13}, SU8/polycarbonate\textsuperscript{14}, polyethylene terephthalate (PET)\textsuperscript{15} and poly (methyl methacrylate) (PMMA)\textsuperscript{16}. Photoluminescence (PL) spectra are recorded at the moment when a mechanical strain is exerted. Similar findings from the aforementioned literature have been reported: a red shift of PL emission energy and decreased peak intensity under uniaxial tensile strain, conforming to the direct-to-indirect transition of the optical band gap as predicated by theoretical studies. The observation can be qualitatively understood as a result of reduced orbital overlap and hybridization due to weakened atomic bonds\textsuperscript{16}. In complementary to the above results, a blue shift of the PL peak and an increase of the emission intensity have been reported for biaxial compressive strain exerted to tri-layer MoS$_2$ through a piezoelectric substrate\textsuperscript{17}. At the device level, back-gated MoS$_2$ transistors on a flexible substrate measured in the stretched state shows a shift of the transfer curve toward lower back-gate voltages and an increase in electron current than the results measured in flat state\textsuperscript{18}.

In this paper, we explore on the concept of “strain-gated” MoS$_2$ MOSFET. Here, strain is exerted to MoS$_2$ channel through the deposition of a silicon nitride stress capping layer that covers the entire transistor active area, analogous to the industry-proven technique applied to the early generation of n-channel silicon transistors. To enhance the field-effect mobility, uniaxial tensile strain along the transistor channel is favored to be generated in order to reduce the band gap and electron effective mass. Current on/off ratio and other performance metrics are measured as the transistors evolve from back-gated, to top-gated and finally, strain-gated structure.

**Results and Discussion**

First, we design sample structure to visualize the internal stress of the silicon nitride indirectly through the Raman mode shift of MoS$_2$ underneath. Silicon nitride (SiN$_x$) deposited by plasma-enhanced chemical vapor deposition (PECVD) are used in both Raman measurement and transistor characterization. As the ions do not respond to the field at a single RF frequency of 13.56 MHz, a thin film deposited under such conditions typically exhibits tensile stress\textsuperscript{19,20}. As shown in Fig. 1(a) and (b), a bi-layer MoS$_2$ sample was exfoliated and transferred.
onto a Si/SiO₂ substrate. The sample is divided into three regions: only region 2 is covered by 125 nm PECVD silicon nitride; the intact regions 1 and 3 serve as control groups. This is to minimize potential non-uniformities from different pristine bi-layer samples.

Before nitride deposition, the bi-layer thickness and the homogeneity of the MoS₂ sample are confirmed by PL and Raman measurement. In all three regions, PL spectra show the two prominent emission peaks at 670 and 627 nm (Fig. 1c–e), corresponding to the two resonances known as A₁ and B₁ excitons, and the Raman spectra show a wavenumber difference of 22 cm⁻¹ between the E₂g and A₁ peak, (Fig. 1f–h), a signature of bi-layer MoS₂. To show the similarities in peak intensity among the three regions, the spectra taken from the as-prepared sample in Fig. 1 are re-plotted in overlapping format in Fig. S1 (Supplementary Information).

In the post-deposition PL spectrum of region 2 (Fig. 1d), we find ~10 nm red shift of the A₁ excitons (from 670 to ~680 nm), and about 40% decrease of the emission intensity. The decrease in peak emission energy is ~27 meV by using E = 1240/λ. The findings indicate a narrowed indirect bandgap possibly due to the tensile stress from the nitride capping layer. Difference in thermal expansion coefficient between MoS₂ and silicon nitride could be another source of stress. As the nitride layer was deposited at 120 °C, a tensile strain can be induced in the MoS₂ flake during the subsequent cooling of the sample to room temperature. For the Raman spectra, given the fact that the referential peak of silicon remains at 520.7 cm⁻¹ before and after nitride deposition, the post-deposition spectrum of region 2 shows a red shift of the A₁g peak (Δω ≈ 1.5 cm⁻¹), while the shift of E₂g peak is negligible (Fig. 1g). We also observed this shift of A₁g peak in other locations on the MoS₂ sample covered by the nitride slab, however no spatial correlations are concluded. The shift of the A₁g peak also disappears once the laser crosses either left or right boundary of the nitride slab. This change in Raman spectra could be due to an extra strain in the out-of-plane direction, but we are still trying to identify the reason behind. Regarding Region 1 and 3, no noticeable peak shift was observed either in PL or Raman characterization, the slightly dropped intensity might originate from some resist residue adsorbed on the MoS₂ surface at the lift-off step. The continuously narrowed band gap of bi-layer MoS₂ under tensile strain is calculated and shown in Fig. 2(a). The indirect band gap of bi-layer MoS₂ at 0% strain is 1.25 eV. It is clear that the band gap decreases as the tensile strain increases, the tunability on average is ~100 meV per 1% tensile strain. Beyond 10%, the bi-layer MoS₂ finally becomes close to metallic. Based on the simulated result, the magnitude of the tensile strain induced by the nitride capping layer is estimated to be much less than 1%. From the study of band structure in Fig. 2(b–d), the decrease in band gap is mainly coming from the lowering of the conduction band at K point. The band structure keeps indirect through all different tensile strain conditions. And the contribution of the band edge comes purely from Mo element’s d-orbital.

Next, we design experiments to apply the nitride stress capping layer onto MoS₂ field-effect transistors and observe the strain's influence on device performance. Figure 3 shows the three main stages for electrical characterization, denoted as back-gate (BG), top-gate (TG) and strain-gate (SG). The key challenge for achieving the final strain-gate configuration is fabricating MoS₂ FETs with robust electrical contacts to endure multiple fabrication and electrical measurement cycles. From our experience, the yield of working FET and the performance repeatability are significantly improved with one-dimensional edge contacts, in contrast to conventional surface contacts which are affected by inherent surface defects found on natural MoS₂ crystal.

Figure 4 shows the optical images captured at different stages of transistor fabrication. Firstly, MoS₂ flakes were directly exfoliated onto a Si/SiO₂ substrate and a bi-layer region (faint purple color) was identified (Fig. 4a). The thickness is confirmed by the 22 cm⁻¹ difference between the E₂g and A₁g modes and low wavenumber shear mode E₂g (20 ~25 cm⁻¹) (Fig. S2). After the first e-beam lithography step, the sites for source drain contacts were patterned and etched by helium plasma (Fig. 4b). Instead of using SF₆ gas, helium plasma can prevent potential oxidation of the exposed MoS₂ atoms. The edge contacts are completed after metal evaporation (Sc/Ni) and lift-off (Fig. 4c). At the completion of the electrical measurement on the back-gate transistors, the bi-layer region is covered by a top-gate dielectric comprised of an AlOₓ/HfO₂ stack (Fig. 4d). The AlOₓ layer, which helps to increase the edge contacts which are affected by inherent surface defects found on natural MoS₂ crystal.

The effective modulation of the drain current by the applied voltage is denoted as back-gate (BG), top-gate (TG) and strain-gate (SG). The key challenge for achieving the final strain-gate configuration is fabricating MoS₂ FETs with robust electrical contacts to endure multiple fabrication and electrical measurement cycles. From our experience, the yield of working FET and the performance repeatability are significantly improved with one-dimensional edge contacts, in contrast to conventional surface contacts which are affected by inherent surface defects found on natural MoS₂ crystal.

In total, three FETs are fabricated on this bi-layer MoS₂, as labeled on Fig. 4e. To test whether all three are functioning, the output and transfer characteristics were initially taken at the back-gate configuration (Fig. S3). The effective modulation of the drain current by V_DS tells that all three edge-contacted FETs are working properly, though some non-linearity and discrepancy is still observed in the output plots, which could be due to the difference in contact resistance. In order to make a fair comparison across the three stages of fabrication, from this point on, the discussion is focused on the electrical performance of FET No. 2.

The full electrical characteristics of FET No. 2, comprising of results from back-gated, top-gated and strain-gated transistor structure, are given in Fig. 5. Output plots are shown in the first row, in which V_DS is swept from −1 to 1 V. The inset of Fig. 5(a) enlarges the area close to the origin. From −200 to 200 mV, I_D changes almost linearly with V_DS, indicating an Ohmic-like junction. From Fig. 5(a) to (c), qualitatively, the spreading of the output curves along the vertical axis indicates continuously improved drain current and a more effective current modulation from the top-gated voltages. Of note is that the top-gate voltages were set to 10 times smaller than that for back-gate, for a much thinner top-gate dielectric and the enhanced over 60 times capacitance per unit area. As shown in the following equation, here we use 30 nm thick HfO₂ with a theoretical relative permittivity of ε_r, = 25, and 285 nm thick SiO₂ with ε_r, = 3.9.


Figure 2. Strained bi-layer MoS₂ band structure calculation. (a) The continuously narrowed band gap of bilayer MoS₂ as a function of increasing lattice constant, the band structure with 1%, 5% and 10% tensile strain are shown in (b, c and d). The blue and green dots in the figure represent Mo and S element occupation respectively.

\[
\frac{C_{HfO_2}}{C_{SiO_2}} = \frac{\varepsilon_{HfO_2}}{d_{HfO_2}} \times \frac{\varepsilon_{SiO_2}}{d_{SiO_2}} = \left[ \frac{25}{30 \text{ nm}} \right] \times \left[ \frac{3.9}{285 \text{ nm}} \right] \approx 61
\]

(1)

Gate transfer plots are in the second row of Fig. 5. We notice three important results across all three transfer plots: 1) the drain current increases significantly at positive gate voltage, indicating n-type FET; 2) steep slopes of the transfer curves when we switch from back-gate to top-gate; owing to the much improved top-gate capacitance with HfO₂, the change in slope is later manifested as an evident decrease in sub-threshold swing; 3) the continued increase in drain current from top-gate to strain-gate, which proves that the deposition of the nitride stress liner indeed helps to enhance the transistor’s electrical performance.

The extracted carrier transport parameters are plotted in Fig. 6. The “on” and “off” currents are defined as the maximum and the minimum currents in the measured gate voltage range. The mobility values were obtained firstly by curve fitting to each \( I_D - V_{GS} \) curve and then calculated by applying the following equation:

\[
\mu_{FE} = g_m \cdot \frac{L}{W} \cdot \left( \frac{1}{C_{ox}} \right) \cdot \left( \frac{1}{V_{DS}} \right)
\]

(2)

where \( C_{ox} = 1.2 \times 10^{-8} \text{F/cm}^2 \) is the capacitance per unit area of the 285 nm-thick SiO₂ back-gate dielectric, and \( C_{ox} = 7.4 \times 10^{-7} \text{F/cm}^2 \) for 30 nm HfO₂ top-gate dielectric. The aspect ratio of transistor No. 2 is \( W/L = 7.5/1.8 \mu \text{m} \). Note that the maximum value of the transconductance \( g_m = dI_D/dV_{GS} \) was used in the calculation. The data plotted in Fig. 6 were all obtained at \( V_{DS} = 1 \text{ V} \). From back- to top-gated FET, quantitatively, we notice a 5 times
increase in on-current, ~98% decrease in off-current, significantly improved current on/off ratio from <100 to ~10^4, and a 43% increase in electron mobility. For the threshold voltage calculated using the linear extrapolation method, the value shrinks from −14 to −0.8 V; so does the sub-threshold swing, drops from 30 to 1.2 V. These improvements are most likely originated from the higher gate capacitance of the more insulative HfO_2 in the top-gate structure. From top- to strain-gated FET, we find another 58% increase in electron mobility and 46% increase in on-current magnitude, $V_{th}$ decreases slightly to −0.7 V, and the sub-threshold swing remains relatively small at 1.4 V. The consistancy in sub-threshold swing indicates that there is no change of the interface quality of the dielectric layer. Distinct piezotronic response of the MoS_2 was not observed, which is probably because of the opposite orientation of alternating layers in bi-layer 2H MoS_2. Flakes with even number of layers are expected to be centrosymmetric and non-piezoelectric. As the nitride PECVD deposition temperature is lower than that for contact annealing, we can conclude that the performance improvement is not due to a better contact, but a direct result from the strain effect induced by the nitride capping layer. Though the magnitude of the mobility is still too low for practical logic device, which is probably a result of the long span of time in multiple e-beam lithography steps and repetitive electrical characterizations in between, the ideas presented in this paper can be undoubtedly expanded to other TMDC semiconductor with high intrinsic mobilities. FET No. 3 exhibits similar carrier transport enhancement as FET No. 2, the extracted electron mobilities at the corresponding BG – TG – SG are 0.11,
0.13 and 0.20 cm²/Vs, whereas no obvious performance improvement was observed in FET No. 1. It is possible that PECVD gives rough granulated SiNx surface at lower deposition temperature³¹, leading to incoherent stress experienced by the bi-layer MoS₂ flake.

In summary, we have characterized the strain induction in MoS₂ through both spectroscopic study and electrical measurement. We have also seen significant improvement in on-current density and mobility, when the FET evolves from back-gated, to top-gated and finally to strain-gated configuration. We conclude that the pioneering approach of strain induction through a nitride stress liner has a benign effect on improving the carrier transport property. It is also an industry-compatible and permanent solution for strain exertion without relying on external substrate bending facilities. Further studies of the strain effect in “all 2D” transistors with perceived greater flexibility would bring in deeper understanding of the tunability of TMDCs’ electrical properties under mechanical strain.

Figure 4. The process flow for making strain-gated FETs (a) an optical image of a bi-layer MoS₂ in faint purple color on a Si/SiO₂ substrate; (b) the material exposed in source drain sites have been etched by helium plasma; (c) edge contact metal (Sc/Ni, 50/30 nm) evaporation and lift-off; (d) 30 nm HfO₂, ALD deposition (with 2 nm Al seed) and lift-off, for the bi-layer region only; (e) top-gate metal (Ti/Au, 10/60 nm) deposition, the FETs are labeled No. 1, 2 and 3; (f) 125 nm SiNx PECVD deposition at 120 °C, capping both the multi-layer and the device area.
Methods

Bi-layer MoS2 sample preparation. The bi-layer MoS2 sample for spectroscopic study was initially exfoliated onto a PDMS-based gel-film, and then transferred onto a pre-cleaned oxidized Si substrate through an all-dry micro-manipulation procedure. This additional exfoliation step on gel-film significantly improved the yield of mono- and bi-layer MoS2 to almost 100%. The freshly cleaved ultra-thin regions usually come in area of a few μm², feasible for multiple transistor fabrication. The second bi-layer MoS2 sample used in strain-gated transistor fabrication was prepared following the conventional Scotch-tape based mechanical exfoliation using bulk crystal.

Raman and photoluminescence spectra were acquired under ambient conditions with a Horiba LabRAM HR spectrometer equipped with a 523 nm laser supply and an 1800 lines/mm grating. A 100× objective was used for focusing the laser to an approximately 1 μm spot onto the sample. The laser power is <1 mW to prevent sample heating. Thin film thickness measurement was carried out under tapping mode from the AFM module of the spectrometer.

Strained bi-layer MoS2 band structure calculation. The band dispersions and band gap energies of bilayer MoS2 with and without strains were calculated using the Vienna ab initio simulation package (VASP) in the projected-augmented-wave method. The generalized gradient approximation (GGA) of the Perdew-Burke-Ernzerhof form is used for the exchange correlation energy. The vdW interactions between the two MoS2 layers are accounted for by using the DFT-D2 method of Grimme. The kinetic energy cutoff for our calculation is set as 500 eV. The lattice constant of bilayer MoS2 without strain is about 3.179 Å. For all structural relaxations, the convergence tolerance on the Hellmann-Feynman forces is less than 0.01 eV/Å. An 8×8×1 Monkhorst-Pack k-point mesh is used for 2D films. The vacuum layer added to the system is nearly 20 Å. The other strained systems are all calculated under the same setting.

Edge contacts and MoS2 transistors fabrication. A degenerately boron-doped (0.001~0.005 Ω•cm) silicon substrate with 285 nm SiO2 capping layer, served as the global back gate and gate dielectric; silicon nitride PECVD deposition: Plasma-therm Unaxis 790, deposition pressure 900 mTorr, power 25 W, RF frequency...
13.56 MHz, 2% SiN 4 200 sccm, NH 3 4 sccm, N 2 900 sccm, deposition time 10 mins at 120 °C, film thickness 125 nm; helium plasma for opening edge contacts: STS Reactive Ion Etcher Dielectric System, helium flow rate 50 sccm, 75 W, 20 s; HfO2 ALD deposition: Cambridge Nanotech Savannah 100, 250 cycles at 120 °C, film thickness 30 nm; Metal evaporation: Temescal BJD 1800 system, aluminum seed (2 nm, 0.1 Å/s), source drain electrodes (Sc/Ni, 50/30 nm), back-gate electrode (Ti/Au, 10/100 nm), top-gate electrode (Ti/Au, 10/60 nm); post-HfO2 deposition contacts annealing: RTP-600S system, 200 °C, Ar, 1 hour. The FETs were characterized in a shielded probe station connected to an Agilent 4155 C semiconductor parameter analyzer. The entire measurement was carried out at room temperature in air.

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