Article

Launch Bar Dynamics Character Analysis of Carrier-Based Aircraft Catapult Launch

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Abstract: The launch bar is a unique device of carrier-based aircraft, which is connected to the nose gear and shuttle. In order to avoid the launch bar striking the flight deck after the launch bar pops out of the shuttle, it is very important to research the dynamics performance of the launch bar. This paper establishes a staged mathematical model of catapult launch including the launch bar, a steam-powered catapult, a holdback bar, and a carrier-based aircraft. This article studied the effect of the mass of the launch bar, restoring moment of the launch bar, and center of gravity position of the launch bar on the dynamics performance of the launch bar. The results showed the following: (1) we could reduce the risk collision of the launch bar and deck by reducing the mass of the launch bar, increasing the restoring moment, and shifting the center of gravity position of the launch bar; (2) under the working condition of this article, we changed the center of gravity position of the launch bar to control the sink of the launch bar end, having the most obvious effect, and we reduced the mass of the launch bar, having the least effect on controlling the sink of the launch bar end; however, reducing the mass of the launch bar could also greatly reduce the risk collision of the launch bar and deck; (3) in order to avoid the launch bar striking the flight deck, the restoring moment of the launch bar must overcome the sum of other moments. The study results can give a theoretical reference for designing and testing the launch bars of carrier-based aircraft. It can also give a theoretical reference for designing and testing the launch bar’s driving mechanisms.

Keywords: catapult launch; launch bar; holdback bar; steam catapult; carrier-based aircraft

1. Introduction

The main role of the launch bar is to transmit the catapult force [1–3] to the aircraft, allowing the aircraft to accelerate while taxiing before reaching a safe catapult end airspeed within the limited length deck. At the end of the catapult [4–6], the launch bar pops out of the shuttle. The launch bar, which is in the air, needs a mechanism for controlling and overcoming other damping torque. The launch bar may strike the flight deck if the restoring moment of the launch bar cannot overcome the other damping torque. Therefore, restraining the contrarotation of the launch bar is an important condition when ensuring the secure catapult launch of an aircraft. It is very important to research the characteristics of launch bar dynamics.

When researching the launch bar for a catapult launch, Reference [7] indicated that the launching system should consist of launch bar installation, cockpit controls and components, and holdback and release installation. Small [8] studied full-scale tests of nose tow catapulting, and showed that nose wheel tow catapulting was better than a bridle launching system. Reference [9] established a six-degree-of-freedom dynamics model of a catapult launch for carrier-based aircraft taking into consideration aircraft off-center position, and the launch bar load was calculated. Reference [10] carried out parametric models and optimization analyses of launch bar driving mechanisms. There were
some patents introduced concerning the installation of launch bars, for example, (a) an improved power-operated launching system for aircraft [11]; (b) launching and holdback gear designed around the nose landing gear of the aircraft [12]; (c) a launch bar pivotally mounted on the oleo strut torque links [13]; (d) a control mechanism for controlling a catapult bar [14]; and (e) a nose gear assembly using one nose wheel for a catapult-launched airplane [15]. However, there is no literature on the characteristics of launch bar dynamics after the launch bar automatically disengages from the shuttle at the end of the power stroke.

On the basis of the security of a catapult launch for carrier-based aircraft, Reference [8] indicated that nose wheel tow would increase the safety of the flight deck. Lucas [16] researched catapult criteria for carrier-based airplanes. Many studies were carried out on the problem of the main gear’s off-center location for the catapult launch of carrier-based aircraft [7,9,17–20]. Reference [21] indicated that the catapult minimum end airspeed was the airspeed value that could be safely tested aboard an aircraft carrier. In Reference [22], a total of 39 launching tests were performed during the first test campaign; the conclusion was that the launch bar retraction system had to be improved. Further research on catapult launch mainly focused on the effect of deck motion [23–25], the control problems of the climbing stage [26–29], the effect of nose landing gear’s sudden extension [30–32], the analysis of aircraft–carrier parameter matching [33,34], and so on. However, despite all public documents in regard to catapult launch around the world, there is a lack of research on the effect of the characteristics of launch bar dynamics on the security of the catapult launch after the launch bar automatically disengages from the shuttle at the end of the power stroke.

This paper made great efforts toward the goal of analyzing the characteristics of launch bar dynamics. This article built a mathematical model of catapult launch, including a launch bar dynamics model, and analyzed the effect of some factors on the characteristics of launch bar dynamics, such as the mass of the launch bar, the center of gravity position of the launch bar, and the restoring moment of the launch bar. The main contributions of this paper are as follows: (1) a complete mathematical model of catapult launch used for analyzing the characteristics of launch bar dynamics character was established; (2) the effects of some factors on the characteristics of launch bar dynamics were analyzed, providing a theoretical reference for the relevant design of launch bars.

The rest of this article is organized as follows: A mathematical model of catapult launch is introduced in Section 2. Influence factor analysis of launch bar dynamics is carried out in Section 3. Finally, conclusions are presented in Section 4.

2. Mathematical Model of Catapult Launch

A steam catapult launch system for carrier-based aircraft is made up of a steam-powered catapult, a towing holdback device, and the carrier-based aircraft, and its simplified model is presented in Figure 1. The distinct parameters of the catapult launch process of carrier-based aircraft can be found in Reference [35].

![Figure 1. The general steam-powered catapult launch system for carrier-based aircraft.](image-url)
2.1. Modeling Assumptions

In order to research the basic law of motion for a launch bar during the process of a catapult launch, this article made the following assumptions:

1. The thermodynamic process of the steam catapult is an adiabatic process; the steam in the catapult is not affected by its environment.
2. The accumulator, cylinder, and piston are rigid bodies, neglecting any changes in shape caused by temperature and pressure.
3. The thermodynamic processes of the accumulator and cylinder are quasi-static processes.
4. The steam has no friction with the pipe wall during the progress from accumulator to cylinder.
5. The fuselage of the carrier-based aircraft is a rigid body.
6. The deck motion and airflow interference are neglected.
7. The carrier-based aircraft has no yaw angle during the process of catapult launch.
8. The effects of asymmetric factors are not considered during the process of catapult launch.
9. The nose gear is vertical with the flight deck surface of the carrier during the process of catapult launch.
10. The elevator angle is fixed during the whole process of catapult launch.
11. The force and damper model of the nose gear adopts the classical two-mass spring–damper model which divides the aircraft into two parts: the elastic support mass [36] and the inelastic support mass.

2.2. Mathematical Model of the Launch Bar

2.2.1. The Mathematical Model of the Launch Bar before the Launch Bar Automatically Disengages

The launch bar is connected with the nose gear and shuttle before the launch bar automatically disengages, as presented in Figure 2.

![Figure 2](image_url)

**Figure 2.** The engagement system of the launch bar, nose gear, and shuttle.

In Figure 2, $Sx_sy_sz_s$ is the carrier-body coordinate system which is fixed with the carrier. The origin $S$ is placed at the center of gravity of the carrier. Axis $x_s$ is parallel with the deck surface of the carrier and points to the prow. Axis $z_s$ is in the carrier symmetry plane and points to the bottom, and axis $y_s$ is perpendicular to the carrier symmetry plane and points to the starboard. The point $O$ is the pivot...
point of the launch bar. \( O_{x_0y_0z_0} \) is the pivot point coordinate system which is fixed with the nose gear. Axis \( x_0 \) is perpendicular to the nose gear and points to the front. Axis \( z_0 \) is parallel with the nose gear and points downward, and axis \( y_0 \) follows the right-hand rule. Point \( O_n \) is the center of gravity of the tire of the nose gear. Point \( D_n \) is the contact point of the deck and the tire of the nose gear.

According to the geometrical relationship between the launch bar and nose gear, we can get the equation of the center of gravity (cg) position of the launch bar as follows:

\[
\begin{align*}
    x_{cg} &= x_O + K_{launch} L_{launch} \cos \theta_c \\
    z_{cg} &= z_O + K_{launch} L_{launch} \sin \theta_c
\end{align*}
\]

where \((x_O, z_O)\) is the coordinate of the pivot point, \( K_{launch} \) is a proportionality factor which is defined as the distance from the cg position of the launch bar to the pivot point divided by the length of the launch bar, \( L_{launch} \) is the length of the launch bar, and \( \theta_c \) is the angle between the launch bar axis and the deck.

According to the geometrical relationship between the launch bar and nose gear, we can get the equation of the end of the launch bar as follows:

\[
\begin{align*}
    x_{end} &= x_O + L_{launch} \cos \theta_c \\
    z_{end} &= z_O + L_{launch} \sin \theta_c
\end{align*}
\]

The steam catapult transmits the load to the nose gear via the launch bar; thus, the axial load of the launch bar can be divided into two parts: the catapult force \( F_{launch1} \) which is parallel with the deck, and the force \( F_{launch1} \) which is vertical with respect to the deck; they are given by

\[
\begin{align*}
    F_{launch1} &= F_{launch} \cos \theta_c \\
    F_{launch2} &= F_{launch} \tan \theta_c
\end{align*}
\]

where \( F_{launch} \) is the axial load of the launch bar.

2.2.2. The Mathematical Model of the Launch Bar after the Launch Bar Automatically Disengages

At the end of the catapult, the launch bar pops out of the shuttle, as presented in Figure 3.

![Figure 3. The system of the launch bar, nose gear, and shuttle.](image-url)
The equation of the cg position of the launch bar is as follows:

\[
\begin{align*}
\dot{x}_{cg} &= x_O + K_{\text{launch}}L_{\text{launch}} \sin \theta_r \\
\dot{z}_{cg} &= z_O + K_{\text{launch}}L_{\text{launch}} \cos \theta_r
\end{align*}
\]

(4)

where \( \theta_r \) is the angle between the launch bar and nose gear.

Because the pivot point coordinate system is parallel to the carrier-body coordinate system, the equation of the momentum theorem for the launch bar is given by

\[
M_r - M_g - M_a - M_t = J \ddot{\theta}_r
\]

(5)

where \( M_r \) is the restoring moment of the launch bar, \( M_g \) is the moment of gravity of the launch bar, \( M_a \) is the moment of the force of inertia due to the longitudinal acceleration of the pivot point, \( M_t \) is the moment of the force of inertia due to the nose landing gear’s sudden extension, \( J \) is the moment of inertia of the launch bar, and \( \ddot{\theta}_r \) is the angular acceleration of the launch bar.

The equation of the angle between the launch bar and nose gear is as follows:

\[
\theta_r = \theta_{r0} + \Delta \theta_r
\]

(6)

where \( \theta_{r0} \) is the initial angle between the launch bar and nose gear, and \( \Delta \theta_r \) is the change value of the angle between the launch bar and nose gear.

The equation of the restoring moment of the launch bar is as follows:

\[
M_r = k_r
\]

(7)

where \( k_r \) is the restoring moment of the launch bar, which is a constant in this article.

The equation of the moment of gravity of the launch bar is as follows:

\[
M_g = m_{\text{launch}}gK_{\text{launch}}L_{\text{launch}} \sin \theta_r
\]

(8)

where \( m_{\text{launch}} \) is the mass of the launch bar, and \( g \) is the acceleration due to gravity.

The equation of the moment of the force of inertia due to the longitudinal acceleration of the pivot point is as follows:

\[
M_a = m_{\text{launch}}aK_{\text{launch}}L_{\text{launch}} \cos \theta_r
\]

(9)

where \( a \) is longitudinal acceleration of the aircraft.

The equation of the moment of the force of inertia due to the nose landing gear’s sudden extension is as follows:

\[
M_t = m_{\text{launch}}a_2K_{\text{launch}}L_{\text{launch}} \sin \theta_r
\]

(10)

where \( a_2 \) is the vertical acceleration of the elastic support mass.

The vertical acceleration of the elastic support mass can be calculated as follows:

\[
a_2 = \begin{cases} 
\frac{m_2g-F_n-f_u-F_h-Y_n}{m_2} & \ddot{u} \geq 0 \\
\frac{m_2g+F_u-F_h-Y_n}{m_2} & \ddot{u} < 0 
\end{cases}
\]

(11)

where \( m_2 \) is the elastic support mass of the nose gear, \( F_n \) is the air spring force of the nose gear, \( f_u \) is the strut friction force of the nose gear, \( F_h \) is the hydraulic damping force of the nose gear, \( Y_n \) is the equivalent aerodynamic lift acting on the nose gear, and \( \ddot{u} \) is the gear stroking velocity of the nose gear.

The air spring force can be calculated as follows:

\[
F_n = A_g \left[ p_0 \left( \frac{V_0}{V_0 - A_g \ddot{u}} \right) - p_{atm} \right]
\]

(12)
where \( A_d \) is the effective area of pressure in the air chamber, \( P_0 \) is the initial pressure of air in the chamber, \( V_0 \) is the initial volume of air in the chamber, \( u \) is the stroke of the damper, \( \gamma \) is gas polytropic index which varies between 1.1 and 1.4, and \( P_{atm} \) is the air pressure.

The equation for calculating the hydraulic damping force is shown below.

\[
F_h = \begin{cases} 
0 & \text{if } u = 0 \\
\frac{\rho}{2a} \left( \frac{A_1^2}{A_2^2} u^2 + \frac{A_3^2}{A_4^2} u^2 \right) & \text{if } u \neq 0
\end{cases}
\]

(13)

where \( \rho \) is the density of the oil, \( C_d \) is the coefficient of contraction of the main oil chamber, \( C_{ds} \) is the coefficient of contraction of the back oil chamber, \( A_b \) is the area of oil pressure of the main oil chamber, \( A_d \) is the oil hole area of the main oil chamber, and \( A_s \) is the oil hole area of the back oil chamber.

The strut friction force of the nose gear is displayed below.

\[
f_3 = \frac{\dot{u}[K_m F_a + \mu_b(N_u + N_l)]}{|\dot{u}|},
\]

(14)

where \( K_m \) is the friction coefficient of the leather cup, \( \mu_b \) is the bending friction coefficient of the damper, \( N_u \) is the normal force of the upper supporting points when the damper bends, and \( N_l \) is the normal force of the lower supporting points when the damper bends.

2.3. Mathematical Model of Other Sub-Modules

2.3.1. Mathematical Model of the Steam-Powered Catapult

The steam catapult is mainly made up of an accumulator, a launch valve, a cylinder, and a piston, as shown in Figure 1. Below, the thermodynamics model of the accumulator, the mathematical model of the launch valve, and the thermodynamics model of the cylinder are described. The thermodynamics models of the accumulator and cylinder were verified by comparison with experimental data [3,37].

We can get the differential equation of steam pressure inside the accumulator from Reference [37].

\[
\frac{dP_1'}{dt} = -\kappa R_g T_1' \frac{Q_{m1}}{V},
\]

(15)

where \( \kappa \) is the ratio of specific heat, \( R_g \) is the gas constant of steam, \( Q_{m1} = dM_s / dt \) is the mass flow rate, \( dM_s \) is the outflow mass of steam from the accumulator, \( T_1' \) is the temperature of the accumulator, and \( V \) is the volume of the accumulator.

The launch valve is one of the central devices of the steam catapult launch system. The steam catapult launch system can get different catapult energies by adjusting the stroke of the launch valve and the effective launch valve opening area. In order to make the model of the launch valve mostly conform to real conditions, this article uses the test data describing the stroke of a launch valve and the effective launch valve opening area [38]. These test data are illustrated in Figure 4.

The accumulator and cylinder thermodynamic systems are connected with each other through the mass flow rate. The formula of mass flow rate is given by

\[
Q_{m1} = \begin{cases} 
A_d \frac{P_1^r}{V_{t1}'} \sqrt{\frac{2 \kappa - 1}{\kappa \gamma - 1} \left( \frac{P_1}{P_1'} \right)^{\frac{\kappa - 1}{\kappa}} - \left( \frac{P_1}{P_1'} \right)^{\frac{\kappa - 1}{\gamma - 1}} - \left( \frac{P_1}{P_1'} \right)^{\frac{\kappa - 1}{\kappa \gamma - 1}}} & P_1 > P_k \\
A_d \frac{P_1^r}{V_{t1}'} \sqrt{\left( \frac{2}{\kappa + 1} \right)^{\frac{\kappa - 1}{\kappa \gamma - 1}} \frac{1}{P_1}} & P_1 \leq P_k
\end{cases}
\]

(16)

where \( A_d \) is the equivalent effective cross-sectional area of the launch valve and the pipe between the accumulator and cylinder, \( P_k \) is the critical pressure, and \( P_1 \) is the pressure of the cylinder.
The critical pressure can be calculated as follows:

$$P_k = P_1 \left( \frac{2}{\kappa + 1} \right)^{\kappa}.$$

The equivalent effective cross-sectional area of the launch valve and the pipe between the accumulator and cylinder is given by the following equation:

$$A_e = \frac{1}{\sqrt{\frac{1}{S_X} + \frac{1}{S_d} + \frac{1}{S_f}}}$$

where $S_X$ is the accumulator outlet area, $S_d$ is the effective launch valve opening area, and $S_f$ is the inlet area of the cylinder.

We can get the differential equation of steam pressure inside the cylinder from Reference [37].

$$\frac{dP_1}{dt} = \frac{\kappa R_1 T_1 Q_{m1}}{A_1(X_{t0} + X)} - \frac{\kappa P_1}{X_{t0} + X} \frac{dX}{dt},$$

where $A_1$ is the area of the piston, $X$ is the displacement of the piston, and $X_{t0}$ is the initial displacement of the piston.

The equation of the steam pressure of the piston is as follows:

$$F_T = 2 \times 10^6 (P_1 A_1 - P_2 A_2),$$

where $A_2$ is the area of the non-steam piston, and $P_2$ is the gas pressure of the non-steam cylinder.

2.3.2. Mathematical Model of the Holdback Bar

The holdback bar acts on three stages: tensioning, aircraft thrust with the holdback, and catapult force build-up with the holdback. The aircraft is in a static balance condition on the deck during the above three stages, because it is affected by the holdback load of the holdback bar. The holdback bar has a minimum release element load. The holdback bar is broken when the axial load of the holdback bar reaches the minimum release load during the catapult force build-up with the holdback.

In Reference [39], the minimum release load was calculated as follows:

$$R_{min} = 1.35 \times \left( \frac{P + f + 0.2G}{\cos \beta} \right),$$

where $P$ is the maximum thrust of the aircraft, $f$ is the tension load, $G$ is the maximum design weight of the aircraft, and $\beta$ is the angle between the holdback axis and the deck at release.
The longitudinal component of the axial load of the holdback bar is as follows:

\[ R_1 = R \cos \beta. \]  \hspace{1cm} (22)

The vertical component of the axial load of the holdback bar is as follows:

\[ R_2 = R_1 \tan \beta. \]  \hspace{1cm} (23)

### 2.4. Staged Mathematical Model of the Catapult Launch

The catapult launch process for carrier-based aircraft is mainly divided into nine stages. In this article, the mathematical model of the catapult launch process corresponds to the final eight stages. The mathematical model of the catapult launch process is composed of a static model of tensioning, a static model of full takeoff power, a mathematical model of catapult force build-up with the holdback, a model of release, a dynamics model of the power stroke, a dynamics model of the free deck run, and a dynamics model of fly-away [35]. The mathematical model of the steam catapult launch process of carrier-based aircraft was verified by comparison with experimental data, and the mathematical model of the minimum release load of the holdback bar was consistent with military specifications [37].

### 3. Results and Discussions

The simulation conditions of the catapult launch are shown in Table 1.

| Simulation Conditions                        | Data                   |
|---------------------------------------------|------------------------|
| The type of aircraft                        | A3                     |
| Mass of the aircraft                        | 31,752 kg              |
| Thrust angle                                | 0°                     |
| Wing area                                   | 72.37 m²               |
| Mean geometric chord                        | 3.55 m                 |
| Nose tire undeflected radius                | 0.4 m                  |
| Main tire undeflected radius                | 0.559 m                |
| Initial steam pressure of accumulator       | 1.38 MPa               |
| The type of steam-powered catapult          | C7                     |
| Mass of piston                              | 3000 kg                |
| Catapult stroke                             | 75.3 m                 |
| Deck edge distance                          | 12 m                   |
| The initial angle between the launch bar axis and deck | 30°             |
| The initial angle between the holdback axis and deck | 25°             |
| The coefficient of friction between the tire and deck surface | 0.17         |
| The coefficient of friction between the piston and cylinder wall | 0.17            |
| Air density                                 | 1.2254 kg/m³           |

#### 3.1. The Effect of the Launch Bar Mass on the Launch Bar Dynamics

This article selected the masses of the launch bar as 10 kg, 20 kg, 30 kg, and 40 kg when the simulation was carried out. The height of the pivot point of the launch bar, the height of the cg position of the launch bar, the height of the end of the launch bar, and the angle between the launch bar and the nose gear with the change in mass of the launch bar are shown in Figure 5.
After the launch bar automatically disengaged from the shuttle at the end of the power stroke, the following values are recorded in Table 2 for the different masses of the launch bar: (1) the peak value and the final value of the height between the pivot point and the deck of the carrier; (2) the peak value and the final value of the height between the center of gravity (cg) position of the launch bar and the deck of the carrier; (3) the peak value and the final value of the height between the end of the launch bar and the deck of the carrier; (4) the final value of the angle between the launch bar and the nose gear.

### Table 2. Results of key parameters with different launch bar masses.

| Masses | 10 kg | 20 kg | 30 kg | 40 kg |
|--------|-------|-------|-------|-------|
| \( h_p^p \) | 0.7693 m | 0.7693 m | 0.7693 m | 0.7693 m |
| \( h_p^f \) | 0.4649 m | 0.4556 m | 0.4464 m | 0.4372 m |
| \( h_c^p \) | 0.1607 m | 0.1420 m | 0.1235 m | 0.1055 m |
| \( h_c^f \) | 0.7414 m | 0.7414 m | 0.7414 m | 0.7414 m |
| \( h_r^e \) | 0.4338 m | 0.4164 m | 0.3988 m | 0.3817 m |
| \( h_r^f \) | 0.1266 m | 0.0917 m | 0.0572 m | 0.0212 m |
| \( \theta_f \) | 63.95° | 62.31° | 60.72° | 59.06° |
In Table 2, \( h^p_O \) is the peak value of the height between the pivot point and the deck of the carrier, \( h^f_O \) is the final value of the height between the pivot point and the deck of the carrier, \( h^p_{cg} \) is the peak value of the height between the center of gravity (cg) position of the launch bar and the deck of the carrier, \( h^f_{cg} \) is the final value of the height between the cg position of the launch bar and the deck of the carrier, \( h^p_{end} \) is the peak value of the height between the end of the launch bar and the deck of the carrier, \( h^f_{end} \) is the final value of the height between the end of the launch bar and the deck of the carrier, and \( \theta^f_r \) is the final value of the angle between the launch bar and the nose gear.

The change of the launch bar mass had no effect on the height from the pivot point to the deck, as shown in Figure 5a. This phenomenon occurred because the pivot point and the nose gear were hinged, and the height from the pivot point to the deck changed along with the change in elastic support mass. Due to the mass of the launch bar being smaller than the elastic support mass, it had little effect on the dynamics of the elastic support mass. Therefore, the changeable position of the pivot point was affected by the condition of the launching system, which was at the end of the catapult, and by the performance of the nose gear after the launch bar automatically disengaged from the shuttle at the end of the power stroke.

With the increase in launch bar mass, some values were reduced after the launch bar automatically disengaged from the shuttle at the end of the power stroke, as shown in Table 2 and Figure 5: (a) the peak value and the final value of the height between the cg position of the launch bar and the deck of the carrier; (b) the peak value and the final value of the height between the end of the launch bar and the deck of the carrier; (c) the final value of the angle between the launch bar and the nose gear. This phenomenon occurred because as the launch bar mass increased, so did the moment of gravity of the launch bar and the moment of the force of inertia due to the longitudinal acceleration of the pivot point. These two moments could affect the equation of the momentum theorem for the launch bar. Finally, the increase of the launch bar mass caused the decrease of the above values.

There was an assumption that the initial height from the end of launch bar to the deck was 5 cm. When the mass of the launch bar was 40 kg, the final value of the height between the end of the launch bar and the deck of the carrier decreased by 7.88 cm, whereby the launch bar would strike the deck surface of the carrier. Therefore, the risk of the launch bar striking the deck surface can be decreased by reducing the mass of the launch bar.

3.2. The Effect of the Restoring Moment of the Launch Bar on the Launch Bar Dynamics

This article selected the restoring moments of the launch bar as 50 N·m, 150 N·m, 250 N·m, and 350 N·m when the simulation was carried out. The height of the pivot point of the launch bar, the height of the cg position of the launch bar, the height of the end of the launch bar, and the angle between the launch bar and the nose gear with the change in restoring moment of the launch bar are shown in Figure 6.

After the launch bar automatically disengaged from the shuttle at the end of the power stroke, the key parameters are recorded in Table 3 for different restoring moments of the launch bar.

| Table 3. Results of key parameters with different restoring moments of the launch bar. |
|----------------|----------------|----------------|----------------|
| Restoring Moments | 50 N·m | 150 N·m | 250 N·m | 350 N·m |
| \( h^p_O \) | 0.7693 m | 0.7693 m | 0.7693 m | 0.7693 m |
| \( h^f_O \) | 0.4372 m | 0.4491 m | 0.4618 m | 0.4752 m |
| \( h^p_{cg} \) | 0.1055 m | 0.1289 m | 0.1551 m | 0.1836 m |
| \( h^f_{cg} \) | 0.7414 m | 0.7414 m | 0.7414 m | 0.7414 m |
| \( h^p_{end} \) | 0.3814 m | 0.4035 m | 0.4261 m | 0.4492 m |
| \( h^f_{end} \) | 0.0212 m | 0.0655 m | 0.1108 m | 0.1570 m |
| \( \theta^f_r \) | 59.04° | 61.13° | 62.23° | 65.33° |
would strike the deck surface of the carrier. Therefore, the risk of the launch bar striking the deck surface can be decreased by increasing the restoring moment of the launch bar.

When the mass of the launch bar was 40 kg, the final value of the height between the end of the launch bar and the deck of the carrier decreased by 7.88 cm, whereby the launch bar would strike the deck, as shown in Figure 6a. This phenomenon occurred because the restoring moment of the launch bar restoring moment had little effect on the dynamics of the elastic support mass.

Upon increasing the restoring moment of the launch bar, the following values increased after the launch bar automatically disengaged from the shuttle at the end of the power stroke: (a) the peak value and the final value of the height between the cg position of the launch bar and the deck of the carrier; (b) the peak value and the final value of the height between the end of the launch bar and the deck of the carrier; (c) the final value of the angle between the launch bar and the nose gear. This phenomenon occurred because as the launch bar mass increased, so did the acceleration of the pivot point. These two moments could affect the equation of the momentum theorem for the launch bar. Finally, the increase of the launch bar mass caused the decrease of the moment of gravity of the launch bar and the moment of the force of inertia due to the longitudinal acceleration of the pivot point. These two moments could affect the equation of the momentum theorem for the launch bar. Therefore, the change in launch bar restoring moment had little effect on the dynamics of the elastic support mass.

The change in launch bar restoring moment had no effect on the height from the pivot point to the deck, as shown in Figure 6a. This phenomenon occurred because the restoring moment of the launch bar was smaller than the load in the catapult launch system. Therefore, the change in launch bar restoring moment had little effect on the dynamics of the elastic support mass.

Upon increasing the restoring moment of the launch bar, the following values increased after the launch bar automatically disengaged from the shuttle at the end of the power stroke: (a) the peak value and the final value of the height between the cg position of the launch bar and the deck of the carrier; (b) the peak value and the final value of the height between the end of the launch bar and the deck of the carrier; (c) the final value of the angle between the launch bar and the nose gear. This phenomenon occurred because the launch bar restoring moment could affect the equation of the momentum theorem for the launch bar, whereby a larger launch bar restoring moment resulted in a larger angular acceleration of the launch bar. The risk of collision between the launch bar and the deck can be reduced by increasing the restoring moment.

There was an assumption that the initial height from the end of the launch bar to the deck was 5 cm. When the restoring moment of the launch bar was 50 N·m, the final value of the height between the end of the launch bar and the deck of the carrier decreased by 7.88 cm, whereby the launch bar would strike the deck surface of the carrier. Therefore, the risk of the launch bar striking the deck surface can be decreased by increasing the restoring moment of the launch bar.
3.3. The Effect of the Launch Bar CG Position on the Launch Bar Dynamics

This article selected the proportionality factor \((K_{\text{launch}})\) of the launch bar as 20%, 40%, 60%, and 80% when the simulation was carried out. The height of the pivot point of the launch bar, the height of the cg position of the launch bar, the height of the end of the launch bar, and the angle between the launch bar and the nose gear with the change in proportionality factor of the launch bar are shown in Figure 7.

Table 4. Results of key parameter with different proportionality factors of launch bar.

| Proportionality Factors | 20%      | 40%      | 60%      | 80%      |
|-------------------------|----------|----------|----------|----------|
|  \(h_p\)               | 0.7693 m | 0.7693 m | 0.7693 m | 0.7693 m |
|  \(h_{cg}\)            | 0.6555 m | 0.5294 m | 0.3912 m | 0.2408 m |
|  \(h_{end}\)           | 0.2035 m | 0.1711 m | 0.1393 m | 0.1085 m |
|  \(f_p\)               | 0.7414 m | 0.7414 m | 0.7414 m | 0.7414 m |
|  \(f_{cg}\)            | 0.6327 m | 0.5006 m | 0.3460 m | 0.1697 m |
|  \(f_{end}\)           | 0.1978 m | 0.1395 m | 0.0824 m | 0.0267 m |
|  \(\theta_f\)          | 67.15°   | 64.54°   | 61.92°   | 59.30°   |

After the launch bar automatically disengaged from the shuttle at the end of the power stroke, the key parameters are recorded in Table 4 for different proportionality factors \((K_{\text{launch}})\) of the launch bar.
The change in proportionality factor of the launch bar had no effect on the height from the pivot point to the deck, as shown in Figure 7a. Due to the mass of launch bar being smaller than the elastic support mass, the change in proportionality factor of the launch bar had little effect on the dynamics of the elastic support mass. Therefore, the change in proportionality factor of the launch bar had no effect on the height from the pivot point to the deck.

With the decrease in proportionality factor of the launch bar, the following values were reduced after the launch bar automatically disengaged from the shuttle at the end of the power stroke: (1) the peak value and the final value of the height between the cg position of the launch bar and the deck of the carrier; (2) the peak value and the final value of the height between the end of the launch bar and the deck of the carrier; (3) the final value of the angle between the launch bar and the nose gear. The proportionality factor of the launch bar could affect the equation of the momentum theorem for the launch bar by affecting the moment of gravity of the launch bar, the moment of the force of inertia due to longitudinal acceleration of the pivot point, and the moment of the force of inertia due to the nose landing gear’s sudden extension. With the decrease in proportionality factor of the launch bar, the three above moments were reduced. Therefore, the risk of the launch bar striking the deck surface can be decreased by reducing the proportionality factor of the launch bar.

There was an assumption that the initial height from the end of the launch bar to the deck was 5 cm. When the proportionality factor of the launch bar was 80%, the final value of the height between the end of the launch bar and the deck of carrier decreased by 7.33 cm, whereby the launch bar would strike the deck surface of the carrier. Therefore, the risk of the launch bar striking the deck surface can be decreased by reducing the proportionality factor of the launch bar.

3.4. Contrastive Analysis of Each Influence Factor

The comparative results of the final height values between the end of the launch bar and the deck surface of the carrier are recorded in Table 5.

|   | Mass Increases | Restoring Moment Decreases | Proportionality Factor Decreases |
|---|----------------|----------------------------|----------------------------------|
| 1 | 0.1266 m       | 0.1570 m                   | 0.1978 m                         |
| 2 | 0.0917 m       | 0.1108 m                   | 0.1395 m                         |
| 3 | 0.0572 m       | 0.0655 m                   | 0.0824 m                         |
| 4 | 0.0212 m       | 0.0212 m                   | 0.0267 m                         |

Under the working conditions of this article, we increased the center of gravity position of the launch bar to control the sink of the launch bar end, which had the most obvious effect, and we reduced the mass of the launch bar, which had the least effect when controlling the sink of the launch bar end. Furthermore, reducing the mass of the launch bar could also greatly reduce the risk of collision between the launch bar and the deck, as shown in Table 5. It is worth noting that the degree of influence of these factors on the results of the final values of the launch bar end may change if the working conditions change.

Each comparative moment of the force of the launch bar is shown in Figure 8 under the following conditions: mass of the launch bar = 40 kg; restoring moment of the launch bar = 350 N·m; proportionality factor of the launch bar = 40%.

We assumed that the restoring moment was a constant in this article. At the moment of the launch bar automatically disengaging from the shuttle at the end of the power stroke, only the restoring moment played a positive role when the launch bar was affected by four moments of force. In the three negative moments, the longitudinal moment of the force of inertia due to the acceleration of the pivot point worked less, whereas the moment of gravity of the launch bar had a large contribution and the moment of the force of inertia due to the nose landing gear’s sudden extension had a large contribution.
The moment of gravity of the launch bar changed less during the process of the free deck run. We can know that the moment of gravity of the launch bar changes with the change in angle between the launch bar and the nose gear according to the equation of the moment of gravity of the launch bar. The moment of gravity of the launch bar changed less because the angle between the launch bar and the nose gear changed less.

In the three negative moments, the longitudinal moment of the force of inertia due to the acceleration of the pivot point worked less. This phenomenon occurred because the longitudinal acceleration of the aircraft was smaller than the acceleration due to gravity and the acceleration of the elastic support mass.

The moment of the force of inertia due to the nose landing gear’s sudden extension changed greatly during the process of the free deck run. Furthermore, the first half of the moment of the force of inertia due to the nose landing gear’s sudden extension played a negative role, and the second half of it played a positive role. The moment of the force of inertia also changed greatly upon the sudden discharge of the holdback load in Figure 8. We should pay attention to the nose landing gear’s sudden extension and the sudden discharge of the holdback load during the process of landing gear design, so as to control launching safety.

In order to avoid the launch bar striking the deck surface, the restoring moment must overcome the total negative moments upon the launch bar automatically disengaging from the shuttle at the end of the power stroke.

4. Conclusions

There is no literature on the characteristics of launch bar dynamics after the launch bar automatically disengages from the shuttle at the end of the power stroke. This paper addressed the problem of the dynamics of the launch bar after the launch bar pops out of the shuttle. In this paper, several points were summarized as follows:

1. A complete mathematical model of catapult launch used to characterize the dynamics of the launch bar was established. Several experimental data from previous research were used to verify the proposed model. The mathematical model of catapult launch including the launch bar can be used to predict the dynamic behavior of the launch bar in order to avoid the launch bar striking the flight deck.

2. The changes in launch bar mass, launch bar restoring moment, and launch bar proportionality factor had no effect on the height from the pivot point to the deck. We can reduce the risk of collision between the launch bar and the deck by reducing the mass of the launch bar, increasing the restoring moment, and shifting the cg position of the launch bar.
3. Under the working conditions of this article, we increased the center of gravity position of the launch bar to control the sink of the launch bar end, which had the most obvious effect, and we reduced the mass of the launch bar, which had the least effect when controlling the sink of the launch bar end. Furthermore, reducing the mass of the launch bar can also greatly reduce the risk of collision between the launch bar and the deck. In order to avoid the risk of collision between the launch bar and the deck after the launch bar automatically disengages from the shuttle at the end of the power stroke, the restoring moment of the launch bar must overcome the sum of the other moments. It is worth noting that the degree of influence of these factors on the results of the final values of the launch bar end may change if the working conditions change.

4. The moment of the force of inertia due to the nose landing gear’s sudden extension changed greatly during the process of the free deck run. Furthermore, the first half of the moment of the force of inertia due to the nose landing gear’s sudden extension played a negative role, and the second half of it played a positive role. The moment of the force of inertia also changed greatly following the sudden discharge of the holdback load. In order to improve launching safety, we should pay attention to the nose landing gear’s sudden extension and the sudden discharge of the holdback load during the process of landing gear design.

5. Selecting the optimal parameters for the launch bar is an important task, so as to control launching safety. The study results can give a theoretical reference for designing and testing the launch bars of carrier-based aircraft. It can also give a theoretical reference for designing and testing the launch bar’s driving mechanisms.

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