Structural, Optical and Magnetic Properties of YIG Doped with Erbium and Terbium Nanofilms using Sol-Gel Method

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Abstract. The combined substitution of Tb and Er ions at the yttrium site with a formula Y_{3-x-y}TbxEr_yFe_{5}O_{12}, (x = 0.0, 0.2) and (y = 0.0, 0.2, 0.6, 1.0 and 1.4) were prepared by a sol-gel method. The crystalline structures of all films were found to have a pure single phase garnet structure. This result confirmed the successful incorporation of erbium and terbium ions into YIG structure. The lattice constant increased linearly with Er^{3+} concentrations from 12.379 to 12.532 Å. The morphology measurement revealed the formation of nanoparticles films. The magnetic properties characterization exhibited that all films are soft ferrimagnetic materials. The results showed the saturation magnetization Ms of pure YIG at 150 (emu / cm³). The Ms value decreased to 110 emu / cm³ at concentration of (x=0.2, y=0.0), then increased with increment of (y) content in a linear manner. The optical transparency of films was over 96 % at (y=0.6) in the wavelength range of 600 to 900 nm. At a short wavelength (λ < 400 nm), absorption coefficient α increased linearly with increment of Er^{3+} ions. Because of nearly zero absorption coefficient in the NIR region, the films could be useful as optical isolator.

Keywords: Nanoparticles, Sol-gel preparation, Er, Tb - substituted YIG, Magnetic properties.

1 Introduction

Yttrium iron garnet (YIG) has low-coercivity, high-remanence, and low conductivity, making them suitable for microwave operation [1-3]. Garnet has complicated structure consists of three crystallographic sites: Dodecahedral, Octahedral and Tetrahedral site [4]. The large dodecahedron sites (2.40Å) are more appropriate for large radii ions such as Bi^{3+}, Y^{3+} or other rare earth ions, whereas the octahedron (2.01 Å) and tetrahedron sites (1.87 Å) are more suitable for a smaller one, such as Al^{3+}.Ga^{3+} and Fe^{3+} [5-7]. In Yttrium iron garnet (Y_{3}Fe_{5}O_{12}), non-magnetic Y^{3+} ions are located on the dodecahedral 24(c) sites while magnetic Fe^{3+} ions occupy the tetrahedral 24(d) sites and the octahedral 16(a) sites in a ratio of 3:2. As the Fe^{3+} ions sit on crystallographically non-equivalent sites, they form two different magnetic sublattices with antiferromagnetic interactions Fe^{3+}_{tet}(↓) → O^{2−} → Fe^{3+}_{oct}(↑) leading to spontaneous magnetization. This interaction will be stronger when the interionic Fe^{3+}-O^{2−} distance are small and the Fe^{3+}_{tet}(↓) → O^{2−} → Fe^{3+}_{oct} bond angle is close to a straight angle (180°) [8, 9].

The garnet structure is compliant with many substitutions which would lead to different application needs. Recently, many reports have been done in substituting Bi^{3+}, Ce^{3+} or other rare earth (R) ions with Y^{3+} and replacement of Fe^{3+} ions by Al^{3+} and Ga^{3+} ions [10-16]. The substitution of terbium ions into garnet crystal at high content is restricted due to the large ionic radius of terbium compared with yttrium ions. It is found that Tb content of less than 0.8 restrains the Y ions to occupy the octahedral sites which may lead to distortion of lattice parameter [17-19]. Therefore, in this study, Tb content is chosen at x=0.2.

Double substitution into yttrium iron garnet (YIG) is an extremely powerful method from a technological applications’ point of view to enhance the physical properties of garnet materials, however the number of studies in this field, is limited [20-23]. This study concentrates on combined substitution of Tb^{3+} and Er^{3+}.
ions at the yttrium site. They are chosen for their high verdet constant, beside that larger values of Bohr magneton (9.59 µB for Tb³⁺ and 9.72 µB for Er³⁺) \[24\]. Therefore, nanoparticles films with the formula of Y₃₋ₓ₋ₓ TbxErₓFe₅O₁₂ (x =0.0, 0.2) and (y =0.0, 0.2, 0.6, 1.0 and 1.4) were prepared. This study is first time reported here for the best of our knowledge to acquire a further comprehension about the effect of double dopant of erbium and terbium into YIG structure on structural, magnetic and optical properties of obtained films. The sol–gel method is chosen as a suitable approach for synthesis, being an inexpensive and producing good-quality garnet film \[25\].

2 Experimental Method

Stoichiometric mixtures of Yttrium (III) Y(NO₃)₃.6H₂O and Iron(III) Fe(NO₃)₃.9H₂O were dissolved in 2mL of 2-methaoxythanol, then stirred by magnetic stirrer at 15 hours. The Y-Fe solution were refluxed at 80 ± 2 °C at 3hr. The terbium (III), Tb(NO₃)₃.5H₂O and erbium (III), Er(NO₃)₃.5H₂O were dissolved separately in 2 ml of acetic acid, then stirred for 3hr before adding gradually to Y-Fe solution. Then, the mixture solutions were refluxed at 80 ± 2 °C for 3hr. The resulting solution was stirred for two days in order to obtain a homogenous gel. The quartz substrate was heated at 75 ± 2 °C for 1 minute before the deposition process to enhance the adhesion of films on substrates. To obtain the film, 20 µL of gel was put on quartz glass and spin coated at 500 ± 3 rpm for 15 sec, then at 3500 ± 5 rpm for 30 sec. The films were annealed at 900° C in the air for 2 hr at a heating rate of 4° C/min to produce crystalline films.

The crystallographic properties were analyzed by X-ray diffractometer (XRD) using CuKα radiation (λ=0.15406 nm). The obtained peaks were compared according to the reference spectrum of pure YIG card number: (01-073-1377). A field emission scanning electron microscopy (FE-SEM; ZEISS Supra 55VP) was used for surface morphology and to measure the film thickness. An Energy dispersive X-ray spectroscopy (EDX) which was coupled with the FE-SEM was used for elements analysis. An (UV-Vis) spectrophotometer (PerkinElmer-Lambda 950) was used to measure the optical properties in the wavelength range of 190 to 1000 nm. A vibrating Sample Magnetometer (VSM) with a maximum magnetic filled of 6,000 Oe was used for in-plane magnetic measurements at room temperature (~29°C).

3 Result and Discussion

3.1 Structure Characterizations

The crystalline structure (see Figure 1) reveals that all films have a pure cubic garnet structure without any secondary phases. The annealing temperature at 900 °C is adequate to obtain the single phases of the cubic garnet structure of these films. The positions of the main peak (420) are slightly shifted to higher diffraction angles as illustrated in Table 1. This shift is evidence of change in lattice constant and confirmed that Tb³⁺ and Er³⁺ ions have been introduced into YIG structure. To calculate the values of the lattice constant (a), Bragg’s equation combined with d-spacing expression of the cubic cell was used \[26\].

\[ a = \left( \frac{\lambda^2 \beta}{4 \sin^2 \theta} \right)^{1/2} \]  

where \( \beta = (h^2+k^2+l^2) \), (h, k, L) are the Miller indices, \( \lambda = 0.15406 \) nm and \( \theta \) is the diffraction angle of X-ray. Bradley & Jay’s extrapolation function is used to minimize the effect of systematic errors \[27, 28\]. Figure 2 shows the dependence of lattice parameter on the content of Er³⁺ ions. The lattice constant (a) increases linearly with the increment of (y) contents. It is attributed to the differences in ionic radii of Y
(1.015 Å), Er (1.00 Å) and Tb (1.04 Å) ions which occupy the dodecahedral sites [29]. Similar results have been reported by Ftema et al. [17] for Tb$_{0.8}$Y$_{2.2}$Al$_{x}$Fe$_{5-y}$O$_{12}$, (0 ≤ y ≤ 1) films. The crystallite sizes were computed using Scherrer’s equation according to the XRD line broadening of the main peak (420).

$$D_{(h, k, l)} = \frac{k\lambda}{\beta \cos \theta}$$

where $D_{(h, k, l)}$ is the average crystallite size in nm, $k$ is the shape factor (0.89), $\lambda$ is the wave length of X-ray, $\beta$ is the half maximum – line width (FWHM) and $\theta$ is diffraction angle. The results are tabulated in Table 1. It changes slightly with the increment of y content.

| Content ($y$) | 2θ (°) at (420) line | Lattice constant (Å) | Crystallite size (nm) |
|--------------|----------------------|----------------------|----------------------|
| 0.0          | 32.415               | 12.379               | 22                   |
| 0.2          | 32.472               | 12.412               | 24                   |
| 0.6          | 32.473               | 12.525               | 28                   |
| 1.0          | 32.474               | 12.532               | 23                   |
| 1.4          | 32.501               | 12.545               | 25                   |

![Figure 1](image1.png) XRD patterns of Y$_{3-x-y}$Tb$_x$Er$_y$Fe$_5$O$_{12}$ films.  
![Figure 2](image2.png) The variation of lattice constant with Er content.

### 3.2 Magnetic Characterizations

The paramagnetic contributions of bare substrates were deducted from total magnetization of deposited films. Figure 3(a) shows the in-plane M-H hysteresis loops of obtained films. The hysteresis loop with normal shape and low coercivity as illustrated in figure 3(b) reveals that all films are soft magnetic materials. The pure YIG ($x=0$, $y=0$) exhibits saturation magnetization, $M_s$ of 150 emu/cm$^3$ which is higher than the $M_s$ of bulk YIG (138.3 emu/cm$^3$) [30] and close to the film prepared using the PLD method on GGG substrate reported by Kehlberger (144±7 KA/m) [12]. As mentioned above, Fe$^{3+}$ ions in a-site and d-site are anti-parallel coupled, leaving a net moment of Fe$^{3+}$ ions in d-site and therefore the $M_s$ value of YIG is given mainly from d-site. On the other hand, in this study some Tb$^{3+}$ & Er$^{3+}$ ions are substituted for Y$^{3+}$ at c-sites at Y$_{3-x-y}$Tb$_x$Er$_y$Fe$_5$O$_{12}$ formula. Both Er$^{3+}$ & Tb$^{3+}$ ions prefer the c- sites due to large ionic radii.
The Tb$^{3+}$ ions doped YIG reduce the saturation magnetization $M_s$ from 150 to 110 emu/cm$^3$ (see Figure 3(a)). This reduction in the $M_s$ value, compared to pure YIG, can be explained by the magnetic moment of the Tb$^{3+}$ ions that coupled antiparallel with the resultant magnetic sublattices of Fe$^{3+}$ ions. Ftema et al. have also reported that doping Y$^{3+}$ with Tb$^{3+}$ reduce the $M_s$ value of YIG film, due to the same reason [17]. Similar behavior had been observed when Er$^{3+}$ ions was doped into YIG, the $M_s$ value is decreased from 150 to 140 emu/cm$^3$ (see Figure 3(a)). Ramadan et al. have reported that the $M_s$ value of Er doped YIG films decreases with the increment of Er$^{3+}$ which due to the Er$^{3+}$ ions that coupled antiferromagnetically with the resultant magnetic sublattices of Fe$^{3+}$ ions [32]. However, when Er$^{3+}$ was doped into Y$_{3-x}$ Tb$_x$ Fe$_5$ O$_{12}$, the $M_s$ increases with increasing of the content (Er) by linear manner (see Figure 4). Nevertheless, our results show that increasing the Er$^{3+}$ in Y$_{3-x}$ Tb$_x$ Fe$_5$ O$_{12}$ increases the $M_s$ values, explaining that the Er$^{3+}$ ion could be coupled ferromagnetically with the resultant magnetic sublattices of Fe$^{3+}$ ions. Hence, the magnetic structure of such mixed garnet could be understood according to the cation distribution as

$$\begin{bmatrix}
 Y^{3+} & Er^{3+} & Tb^{3+} \\
 - & & \\
 Fe^{3+} & & Fe^{3+}
\end{bmatrix}.$$  

The Er$^{3+}$ ions which have a large magnetic moment (9.72 µB) in c-sites align parallel to the net Fe$^{3+}$ ions moment in d-sites [33]. Thus the saturation magnetization increased. The difference in ion coupling has also been reported for Gd$^{3+}$ doped YIG. Aoki et al. have reported that Gd$^{3+}$ could be aligned ferromagnetically with Fe$^{3+}$ ions in the d-site of Ce:Gd-YAIg nanocrystal, resulting increased in $M_s$ values as the Gd$^{3+}$ ion increases for their Ce:Gd-YAIg nanocrystal [34]. However, Xu et al. have reported that Gd$^{3+}$ions align antiferromatically with Fe$^{3+}$ ions in the d-site of Ce:Gd-YIG, resulting decrease in $M_s$ values as the Gd$^{3+}$ ion increases, for their Ce:Gd-YIG samples [23].

The coercivity ($H_c$) for all Er contents (y) exhibited a slight increment and in the range from 21 to 27.5 Oe (see Figure 4). The coercivity are affected by the grain size as shown in Figure 5. The $H_c$ value becomes smaller as the grain size increases. Therefore, the obtained results agree with the theory that the coercivity of small particle ($H_c$) is inversely proportional to grain size according to the relation: $H_c \sim 1/D^\nu$ where $D$ is the particle diameter and $\nu$ is constant [35].

![Figure 3](image_url)

**Figure 3.** (a) The in-plane M-H hysteresis loops of Y$_{3-x}$ Tb$_x$ Er$_y$ Fe$_5$ O$_{12}$ films and (b) shows the hysteresis loop with normal shape and low coercivity.
3.3 Surface Morphology Characterization

All samples were coated with a thin layer of platinum (Pt) using plasma sputtering device to overcome the electrostatic charging during characterization. Figure 6 (a-f) shows the microstructure obtained from FE-SEM at magnification of X50 000. The small white grains on the surface belong to Pt element. All films exhibit the formation of agglomeration which contain nanoparticle. The estimated average grain sizes using the Smart TIFF software are in the range of (29-46 nm). The grain size changes unsystematically with the Er content.

Figure 6. FE-SEM micrographs of Y_{3-x-y}Tb_x Er_y Fe_5 O_{12} films: a (x=0.0, y=0.0), b (x=0.2, y=0.0), c (x=0.2, y=0.2), d (x=0.2, y=0.6), e (x=0.2, y=1.0) and f (x=0.2, y=1.4). The arrow sign shows the voids formation and the circle shows the agglomeration.
The cross section backscattered FE-SEM microstructure was used to measure the film thickness as shown in Figure 7. Even though the preparation parameters were kept the same, the fluctuations of thickness with \((y)\) content can be observed in a non-systematic way which attributed to agglomeration and annealing temperature. The thickness values of obtained films were measured in the range from 122 to 217 nm. However, a good adhesion of the film to the substrate can be observed in Figure 7.

![Figure 7. The typical cross-sectional image of \(\text{Y}_{3-x-y}\) \(\text{Tb}_x\) \(\text{Er}_y\) \(\text{Fe}_5\) \(\text{O}_{12}\) films.](image)

The EDX analysis results of sample with \(y = 0\) and 1.4, are illustrated in Figure 8. All peaks of the elements of Tb, Er, Y, Fe and O appear on the spectrum as well as the Si peak which belongs to the quartz substrate. This result reveals that both the Er and Tb ions are introduced into YIG structure as confirmed by XRD analysis.

![Figure 8. The typical EDX spectroscopy spectrum of \(\text{Y}_{3-x-y}\) \(\text{Tb}_x\) \(\text{Er}_y\) \(\text{Fe}_5\) \(\text{O}_{12}\) films. A\((y=0)\) and B\((y=1.4)\).](image)

### 3.4 Optical properties

The optical transmittance spectra of obtained films at different \((y)\) contents is shown in figure 9. All the samples are highly transparent in the visible and near infrared regions. The \(\text{Tb}^{3+}\) ions doped YIG reduce the optical transparency from 85% to 65%. This reduction, compared to pure YIG, could be attributed to the difference in grain size and thickness. The transmittance improves (above 97%) when \(\text{Er}^{3+}\) ions were doped into \(\text{Y}_{3-x}\) \(\text{Tb}_x\) \(\text{Fe}_5\) \(\text{O}_{12}\) \((y=0.6)\) at near infrared region. The film is the thinnest film (122.8 nm) compared to other films. The slight reduction in transparency of the sample \((y=1.4)\) which has the smallest grain size is mainly due to reflection losses [39]. However, the obtained results are higher than the result reported (90 %)
by Ftema et al. [17]. As a result of the high optical transparency, the surface reflection losses were negligible. One of the main reasons for this high transparency could be attributed to reduction of light scattering from a smaller grain size (y=0.2, 0.6, 1.0 and 1.4), at visible and near infra-red region. The optical absorption coefficient (α) measured from the film’s optical transmission spectra using the Beer – Lambert’s law.

\[ \alpha = -\left( \ln \frac{T}{t} \right) \]  

(1)

where \( T \) is transmittance and \( t \) is film thickness. The change of absorption coefficient against wavelength with different \( y \) content was shown in Figure 10. In general, at a short wavelength (\( \lambda < 400 \) nm), absorption coefficient \( \alpha \) increases linearly with increment of Er\(^{3+}\) ions. This is because probably of electronic transitions of Fe\(^{3+}\) cations in octahedral and tetrahedral sites as well as a charge transfer process between Fe\(^{3+}\) (3d orbit) and O\(^2-\) (2p orbit) ions, which affected by the presence of Er\(^{3+}\) ions [40, 41]. The absorption coefficient becomes lower at visible region, then vanishes in the near infrared, which attributed to crystal field transitions.

![Figure 9. The optical transmission spectra of \( Y_{3-x-y} Tb_x Er_y Fe_5 O_{12} \) films.](image1)

![Figure 10. Absorption coefficient of \( Y_{3-x-y} Tb_x Er_y Fe_5 O_{12} \) films.](image2)

4 Conclusion

The nanoparticles films with the formula of \( Y_{3-x-y} Tb_x Er_y Fe_5 O_{12} \) (\( x = 0.0, 0.2 \)) and (\( 0 \leq y \leq 1.4 \)) were successfully prepared using the sol-gel method. The XRD patterns revealed the cubic structure with single phase of pure garnet. The annealing temperature at 900 °C was adequate temperature to crystalline the films. A slight shift has been observed of the main peaks of XRD pattern to higher diffraction angles due to the incorporation of Er and Tb ions into YIG structure. Generally, higher saturation magnetizations of all films comparing to bulk YIG were reported. These results for such mixed garnet could be understood according to the cation distribution in YIG structure. In terms of hysteresis behavior, the films as a soft materials make them a good candidate for microwave applications. The surface morphology of the obtained films showed an average grain size around (29 – 46 nm) with a homogenous structure. Adding Er\(^{3+}\) ions increases the transparency from 65% to 97% at the visible and near infrared region. The small optical coefficient \( \alpha \) (in cm\(^{-1}\)) in the wavelength region (> 800 nm) enabling the films to be applicable to the optical isolator.
Acknowledgement. This work was financially supported by the Malaysian Ministry of Science and technology under grant No.DLP-2014-006.

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