COSMIC STAR FORMATION HISTORY AND ITS DEPENDENCE ON GALAXY STELLAR MASS

Stéphanie Juneau,1,2 Karl Glazebrook,3 David Crampton,4 Patrick J. McCarthy,4 Sandra Savaglio,3 Roberto Abraham,5 Raymond G. Carlberg,5 Hsiao-Wen Chen,6,7 Damien Le Borgne,5 Ronald O. Marzke,6 Kathy Roth,6 Inger Jorgensen,9 Isobel Hook,10 and Richard Murowinski2

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Abstract

We examine the cosmic star formation rate (SFR) and its dependence on galaxy stellar mass over the redshift range 0.8 < z < 2 using data from the Gemini Deep Deep Survey (GDDS). The SFR in the most massive galaxies (M* > 10^{10.8} M_⊙) was 6 times higher at z = 2 than it is today. It drops steeply from z = 2, reaching the present-day value at z ≈ 1. In contrast, the SFR density of intermediate-mass galaxies (10^{10.2} M_⊙ ≤ M* < 10^{10.8} M_⊙) declines more slowly and may peak or plateau at z ≈ 1.5. We use the characteristic growth time t_{SFR} ≡ ρ_{a∗}/ρ_{SFR} to provide evidence of an associated transition in massive galaxies from a burst to a quiescent star formation mode at z = 2. Intermediate-mass systems transit from burst to quiescent mode at z ~ 1, while the lowest mass objects undergo bursts throughout our redshift range. Our results show unambiguously that the formation era for galaxies was extended and proceeded from high- to low-mass systems. The most massive galaxies formed most of their stars in the first ~3 Gyr of cosmic history. Intermediate-mass objects continued to form their dominant stellar mass for an additional ~2 Gyr, while the lowest mass systems have been forming over the whole cosmic epoch spanned by the GDDS. This view of galaxy formation clearly supports “downsizing” in the SFR where the most massive galaxies form first and galaxy formation proceeds from larger to smaller mass scales.

Subject headings: galaxies: evolution — galaxies: formation

1. INTRODUCTION

The evolution of the global star formation rate (SFR) provides a sensitive probe of galaxy formation and evolution. The earliest determinations of the evolving star formation rate density (SFRD) showed a steep decline from z ~ 1 to the present (Lilly et al. 1996; Madau et al. 1996). The behavior of the SFRD at early epochs (z > 1) remains uncertain owing to variations among SFR diagnostics and poorly constrained (yet potentially large) extinction corrections in the primary rest-frame UV diagnostics (Steidel et al. 1999). Despite these challenges, Hopkins (2004) recently compiled results from 33 studies over the range z = 0 to z ≈ 6. The data were used to constrain the luminosity function of star-forming galaxies and were found to be consistent to within a factor of 3 over 0 < z < 6.

Additional insight into star formation histories (SFHs) may be gained by the consideration of other physical galaxy properties. Stellar mass is arguably the key parameter. The mass in stars provides a measure of the integral of past galaxy stellar mass assembly, which can be coupled with the instantaneous SFR to give a more complete view of galaxy evolution. Mass-based evolution studies are far more deterministic, since unlike luminosity, mass evolution is monotonic. While high-luminosity galaxies often evolve into low-luminosity systems, massive galaxies at early epochs must have descendants among the present massive galaxy population. Recent advances in the modeling of multicolor, and particularly near-IR, selected samples lead to fairly robust determinations of stellar masses for galaxies over a wide range of redshifts and luminosities (Brinchmann & Ellis 2000; Fontana et al. 2004; Glazebrook et al. 2004, hereafter Paper III).

Heavens et al. (2004) inferred the SFH of the universe by modeling the spectra of 96,545 local Sloan Digital Sky Survey (SDSS) galaxies. Their results indicate that SFHs vary strongly with present-day stellar mass. Galaxies with M* < 10^{10.7} M_⊙ go through the peak of their star formation at low redshift (z ≤ 0.5), and M* > 10^{11.2} M_⊙ galaxies show an increasing star formation activity to z > 2. We take advantage of the near-IR selection of the Gemini Deep Deep Survey (GDDS; Abraham et al. 2004, hereafter Paper I) to study directly the evolution of the SFRD as a function of stellar mass at the epoch of observation. Direct measurements of SFRs in the most massive systems at z > 1 were made possible only recently, with the progress of spectroscopy of near-IR surveys probing the redshift range 1 < z < 2. We use (H_α, Ω_m, Ω_Λ) = (70 km s^{-1} Mpc^{-1}, 0.3, 0.7) and Vega magnitudes throughout this Letter.

2. DATA AND SAMPLE SELECTION

The GDDS is a spectroscopic survey of an optical and near-IR selected sample targeting massive galaxies at 0.8 < z <
20. The sample is drawn from the Las Campanas IR imaging survey (McCarthy et al. 2001; Chen et al. 2002) and was designed to select galaxies in the $0.8 < z < 2$ range with an emphasis on the reddest galaxies. Galaxy stellar masses are determined from the mass-to-light ratio based on the GDDS that satisfy (1) (survey limit), (2) a spectroscopic completeness factor that depends on the color and the high-mass galaxies have $M_*$ $\geq 1$0.2 $10^9$ M$_\odot$. The sampling weights are derived as a function of color and magnitude and are used when computing volume-averaged quantities such as the mass density (Paper III) and the SFRD.

The sample selected to compute the SFRD consists of spectra from the GDDS that satisfy (1) $K < 20.6$ (survey limit), (2) a redshift confidence level greater than 75% ($\sigma_{\text{cont}} > 2$ in Paper I notation), and (3) an absence of strong active galactic nuclei activity ($\sigma_{\text{cont}} = 0$ in Paper I notation). From the original GDDS sample of 308 spectra, 211 meet the $K$-detection and redshift confidence class criteria. Strong active galactic nucleus contamination occurs for 1.9% (4/211) of those objects, bringing the final sample to 207 galaxies.

3. METHOD

Given the redshift range spanned by the GDDS, the available SFR indicators are the [O ii] $\lambda$3727 emission line and the luminosity of the rest-frame UV continuum. For the latter, we chose the absolute rest-frame AB magnitude $M_{\text{abs}}$ defined in a synthetic 1900 Å $< \lambda < 2100$ Å box filter using an empirical interpolation scheme from the observed $V$ and $I$ magnitudes (Savaglio et al. 2004). The redshift range is restricted to $z < 1.6$ for SFR([O ii]) measurements, as the [O ii] emission is redshifted out of the optical range at $z > 1.6$. For the SFR($M_{\text{abs}}$) measurements, we use the redshift range where the $M_{\text{abs}}$ interpolation is reliable, i.e., $1.2 < z < 2.0$. Details of the SFR measurements will be given in a forthcoming paper (S. Joneau et al. 2005, in preparation). Briefly, we use the standard SFR($H\alpha$) conversion of Kennicutt (1998) assuming ([O ii]/H$\alpha_{\text{obs}}) = 0.5$ (Glazebrook et al. 1999). Since we use the observed ratio, we need to correct the integrated luminosities at $H\alpha$. We adopt the average extinction of $A_{H\alpha} = 1$ derived from local galaxy samples (Kennicutt 1992).

For SFR measurements based on the rest-frame UV continuum, we apply a dust attenuation correction of the stellar continuum of $A_{V} = 1$. This is the value typically used in the literature when no direct measure of dust obscuration is available (e.g., Lilly et al. 2003; Sullivan et al. 2000). Following the prescription of Calzetti (2001), this correction is $A_{2000} = 2.2$ at 2000 Å. The mean dust obscuration in galaxies depends on the sample selection. K-band selection could include more heavily obscured systems. If the mean extinction is greater by 1 mag, the values of the SFR (§ 4) and the SFRD (§ 5) will shift up by a factor of 2.5 whereas $t_{\text{SFR}}$ (§ 5) will be lower by the same factor.

The SFRD is computed with the $1/V_{\text{max}}$ method and corrected for both sampling and spectroscopic incompleteness. We define spectroscopic completeness factors that depend on the color and the magnitude in the same fashion as the sampling weights from Paper I. We divide the number of spectra with high-confidence redshifts ($\sigma_{\text{cont}} > 2$) by the number of slits in each cell of the color-magnitude plane (see Fig. 12 of Paper I). The median spectroscopic completeness of the sample according to this definition is 82%, consistent with the overall spectroscopic completeness of 79% for the GDDS.

Throughout this Letter, we adopt the Baldry & Glazebrook (2003, hereafter BG03) initial mass function (IMF), which has a very similar slope to the Salpeter (1955) IMF at high masses and provides a good fit to cosmic and galaxy colors locally. The galaxy masses and SFRs based on BG03 can be used interchangeably with those of Salpeter given the conversion $M_{(BG03)} = 0.55M_{(SP)}$; this ratio is virtually independent of SFH to an accuracy of a few percent.

4. STAR FORMATION RATES

The SFRs of the individual galaxies are plotted in Figure 1. The values estimated from the [O ii] luminosity (circles) and those obtained with the continuum luminosity at 2000 Å (triangles) suggest an increase by over 1 order of magnitude in the upper envelope of the SFR values between $z = 0.8$ and $z = 2$. The color of the plotting symbols is keyed to the stellar mass of the galaxies. The mass bin corresponding to the lowest mass galaxies is $10^{9.0} M_{\odot} < M_* < 10^{10.0} M_{\odot}$ (blue), the intermediate mass is defined as $10^{10.0} M_{\odot} < M_* < 10^{10.8} M_{\odot}$ (green), and the high-mass galaxies have $10^{10.8} M_{\odot} < M_* < 10^{11.5} M_{\odot}$ (red). Note that $K$-selected samples will miss contributions from low-mass star-forming galaxies fainter than the $K$ limit. (We show later in Fig. 2 that our full $K < 20.6$ sample underestimates the total SFRD by a factor of 2–3.)

In nearly 10% (20/207) of the galaxies, no [O ii] emission is detected in our spectra. These systems appear to be primarily massive, quiescent galaxies, and their number is uniformly distributed with $z$ (upper limit symbols). The [O ii] flux limit varies from one spectrum to another, depending on the redshift and the integration time for each mask. A conservative flux detect-
 angles represent rest-frame 2000 Å–based rates. The local SFRDs derived from the SDSS by Brinchmann et al. (2004) in our three mass bins are also shown in Figure 2 (squares). These have been transformed to the same IMF and corrected for extinction in the same manner as the GDDS measurements. The comparison with Figure 1 indicates that the increase of SFRD with \( z \) is linked with the increase of SFR in individual galaxies.

Paper III shows that for \( M_∗ > 10^{10.8} M_⊙ \) our sample is mass-complete over our full redshift range. At lower masses, \( 10^{10.2} M_⊙ \leq M_∗ < 10^{10.8} M_⊙ \), we start to become incomplete for \( z > 1.15 \). However, the effect is minor because it is the red objects that start to be missed and these contribute least to the SFRD. In this sense, mass incompleteness also means sensitivity biased toward bluer objects. We can estimate this incompleteness, and calculate a reasonable correction, by bootstrapping from the \( z \leq 1.15 \) complete sample. We calculate that if these objects were placed at \( z = 1.5 \) we would see 90\% of the total SFRD, and at \( z = 1.8 \) we would see 62\%. Note we are effectively assuming that the high-\( z \) sample displays the same range of \( \langle M_{\star} \rangle \) values as the low-\( z \) sample. In practice, we expect the high-\( z \) sample to be bluer, which would make us more complete. Thus, we think our correction is conservative; we extend the upper error bars in Figure 2 to reflect the magnitude of this correction. For \( M_∗ < 10^{10.2} M_⊙ \), we are incomplete throughout our redshift range and plot the SFRD values as lower limits.

The key result of this Letter is illustrated in Figure 2: the cosmic SFRD is a strong function of galaxian stellar mass. The SFRD in our high-mass bin (\( M_∗ > 10^{10.8} M_⊙ \)), while making a minority contribution to the global SFRD, is a factor of \( \sim 6 \) higher at \( z = 2 \) than it is at present. The SFRD of these massive galaxies has strongly declined since \( z = 2 \) and reached the present-day level at \( z \approx 1 \). The SFRD in the intermediate-mass bin (\( 10^{10.2} M_⊙ \leq M_∗ < 10^{10.8} M_⊙ \)) has also steeply declined since \( z \approx 1.2 \) and appears to have reached a peak or a plateau at a redshift of \( \sim 1.5 \). The SFRD in the intermediate-mass bin then declined to the present-day level at \( z < 1 \), and since that time most of the SFRD has been in low-mass galaxies. Our lowest mass bin (\( M_∗ < 10^{10.2} M_⊙ \)) is strongly mass-incomplete. The comparison with the Hopkins points shows that we miss about half the total SFRD at \( z \approx 1 \). This would come from low-mass galaxies fainter than our survey \( K \) limit.

Additional insight into the growth of galaxies can be gained from comparing the stellar mass density to the SFRD over a range of redshifts. The ratio gives a characteristic growth timescale \( t_{\text{sfr}} = \rho_{M}/\rho_{\text{sfr}} \), which can be interpreted simply as the time required for the galaxies to assemble their observed stellar mass assuming that their observed SFR stays constant. Results for the GDDS sample are shown in Figure 3. In this figure, we compare \( t_{\text{sfr}} \), with the Hubble time \( t_H(z) \), the age of the universe at a given redshift. If the values are equal, galaxies can form all their observed stars in a Hubble time. At a given redshift, \( t_{\text{sfr}} > t_H(z) \) suggests that the galaxies are in a declining or quiescent star formation mode at the observed redshift and that the bulk of their star formation has occurred in the past. Conversely, \( t_{\text{sfr}} < t_H(z) \) indicates that the galaxies are going through a burst phase. This past average-to-present SFR allows us to investigate the mode of SFR in the SFRD(\( z \)) diagram. The highest mass galaxies are in a quiescent mode at low-\( z \), transitioning to a burst mode at \( z \approx 1.8 \). The intermediate-mass objects make a similar transition at \( z \approx 1.1 \). The lowest mass systems appear to be observed in a burst mode at all redshifts.
although we note our selection is only sensitive to the bluest objects for this mass-incomplete subsample. The transition redshift from burst mode to quiescent mode is a strong function of stellar mass and looks like a time-delayed echo of the corresponding downturn in the SFRD(z) diagram.

6. DISCUSSION AND SUMMARY

Figures 2 and 3 paint a simple picture of galaxy formation in a volume-averaged sense. The upper envelope of the SFR rises with redshift, and one sees very clear mass-dependent effects in the SFRD(z) and t_{SFR}(z) diagrams. The mass scale of the SFR and of burst activity both decline with time.

The basic conclusion is that the most massive systems formed early and were nearly finished forming their stars by $z \sim 1.5$–2. This is supported by several other lines of evidence. The reddest galaxies in the $1 < z < 2$ interval appear to have ages greater than 1 Gyr and $z_f > 2$ (McCarthy et al. 2004; Cimatti et al. 2004). These objects are $\sim 10^{11} M_\odot$ and suggest that the rising SFRD seen in the high-mass bin of Figure 2 continues to increase, reaching a peak somewhere in the $2 < z < 4$ interval. The most massive local galaxies seen in the SDSS also appear to be dominated by stars with early formation redshifts (Heavens et al. 2004). The Hubble sequence begins to emerge in the range $1 < z < 2$ and appears to be in place at $z \sim 1$ (Brinchmann & Ellis 2000). Disk galaxies that are seen at $z > 2$ are primarily in massive stellar systems (e.g., Labbé et al. 2003; Stockton et al. 2004). As Figure 3 shows, by $z \sim 1$ both the high- and intermediate-mass populations, those that dominated the elliptical and disk portions of the Hubble sequence, have transitioned to fairly quiescent star formation and have formed the bulk of their stars. A puzzle remains in the continued accumulation in the number of massive (see Paper III and Fontana et al. 2004) and/or red-sequence (Bell et al. 2004) galaxies by a factor of $\sim 2$ over $0 < z < 1$. This now has to occur without making a large contribution to the global SFRD in the visible and rest UV, perhaps via mergers that produce either little or heavily obscured star formation.

Our view of galaxy formation reveals a clear “downsizing” of star formation from high-mass to low-mass galaxies with time. This basic picture was first suggested by Cowie et al. (1996). The downsizing picture of galaxy stellar mass assembly is the reverse of the original picture of hierarchical dark matter assembly (Blumenthal et al. 1984), where the large units come last in the sequence. If the hierarchical picture is correct, then star formation must be much more efficient in early times in high-mass systems, as is required (e.g., Paper III) to explain their space densities at high redshift.

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