Proton and Neutron Irradiation Tests of Readout Electronics of the ATLAS Hadronic Endcap Calorimeter

Sven Menke on behalf of the ATLAS Liquid Argon Calorimeter Group

Abstract—The readout electronics of the ATLAS Hadronic Endcap Calorimeter will have to withstand the about ten times larger radiation environment of the future high-luminosity LHC (HL-LHC) compared to their design values. The GaAs ASIC which comprises the heart of the readout electronics has been exposed to neutron and proton radiation with fluences up to ten times the total expected fluences for ten years of running of the HL-LHC. Neutron tests were performed at the NPI in Řež, Czech Republic, where a 36 MeV proton beam is directed on a thick heavy water target to produce neutrons. The proton irradiation was done with 200 MeV protons at the PROSCAN area of the Proton Irradiation Facility at the PSI in Villigen, Switzerland. In-situ measurements of S-parameters in both tests allow the evaluation of frequency dependent performance parameters – like gain and input impedance – as a function of the fluence. The linearity of the ASIC response has been measured directly in the neutron tests with a triangular input pulse of varying amplitude. The performance measurements and expected performance degradations under HL-LHC conditions are presented.

I. INTRODUCTION

THE electronics for the Hadronic Endcap Calorimeter (HEC) [1] of the ATLAS Experiment [2] has been designed to withstand ten years of operation of the LHC at center-of-mass energies of \( \sqrt{s} = 14 \text{ TeV} \) with peak luminosities of \( \mathcal{L} = 10^{34} \text{cm}^{-2} \text{s}^{-1} \) corresponding to an integrated luminosity of \( \int \mathcal{L} = 1000 \text{fb}^{-1} \). The foreseen upgrade to the high-luminosity version of the LHC (HL-LHC) exceeds these limits by up to an order of magnitude. The expected radiation levels for integrated luminosities of 3000 fb\(^{-1}\) in the region of the amplifier electronics located inside the cryostat of the HEC are \( 5.1 \times 10^{13} \text{hcm}^{-2} \) for hadrons above 20 MeV, including safety factors of 5 in total to account for simulation uncertainties, while the corresponding number for 1 MeV equivalent neutrons in silicon is \( 4.1 \times 10^{14} \text{n/cm}^{-2} \), about an order of magnitude larger. The expected total ionization dose (again evaluated for silicon and including the same safety factor of 5) is 6.2 kGy. Proton and neutron irradiation tests have been performed to evaluate the limits of the current readout electronics under HL-LHC conditions.

In addition to the in-situ measurements during the irradiation which took place at room temperature, the neutron irradiated boards were re-measured 3 months after the neutron tests, immersed in liquid nitrogen to simulate the conditions inside the ATLAS endcap cryostat where the readout electronics of the HEC are located.

II. PROTON TESTS

A 198.9 MeV proton beam of 2.71 nA at the Proton Irradiation Facility (PIF) at PSI with a narrow beam profile of \( \sigma_x \approx 7 \text{mm} \) and \( \sigma_y \approx 9 \text{mm} \) at the position of the devices under test has been used to evaluate the radiation hardness against hadrons up to a fluence of \( 2.6 \times 10^{14} \text{p/cm}^{-2} \) after 22.05 h of beam time. A mount with 14 equidistant slots with a slot spacing of 1.7 cm has been placed in the beam line. The 6 boards closest to the beam window housed transistors in SiGe bipolar and Si CMOS FET technology and the GaAs FET based ASICs (BB96) [3]. The BB96 ASICs are currently used in ATLAS for the HEC readout. Slots 7 and 14 were equipped with radiation diodes for dose measurements, and slots 8 – 13 with test structures for the HEC power supplies. A picture of the setup prior to final alignment is shown in Figure 1.

The beam position and width has been measured with sheets of EBT2 radiation film placed for 5 minutes in the beam at 6 different positions. The total beam flux was measured and calibrated by the PIF group with ionization chambers. The raw
Fig. 2. Horizontal (top) and vertical (bottom) profiles of the proton beam at PIF behind slot 3 normalized to the total fluence after 22 hours of beam time. The data points indicate the measurements from the radiation film. Superimposed are Gaussian fits. The yellow bands depict the position of the ASIC and the vertical lines the center and boundaries of the housing. The film was scanned with a resolution of 600 dots per inch (dpi). The Gaussian parameters for mean and width are also given in dpi.

Fig. 3. Picture of the test setup at NPI. The neutrons leave the D$_2$O target at the the upper left part of the picture and travel to the right. The test boards for the HEC readout electronics occupy the first 6 slots and slots 8 and 9 of the aluminum frame in the picture center.

III. NEUTRON TESTS

Neutron tests have been performed with a $12 - 14 \mu A$ and 36 MeV proton beam of the variable-energy cyclotron U-120 at NPI in Rež hitting a 16 mm thick D$_2$O target to produce a neutron spectrum with $E_{\text{kin}} < 32$ MeV and $\langle E_{\text{kin}} \rangle \approx 14$ MeV [4]. Similar to the setup in the proton test, 16 equidistant slots with distances of 1.7 cm between adjacent board centers have been equipped and placed in front of the 3 mm steel plate enclosing the target. The alignment was done prior to mounting the target and the test boards by burning small holes with the proton beam in two calibration boards placed instead of the test structures in the first and last slot. Slots 1–6, 8, and 9 were equipped with the BB96 GaAs ASICs, slots 7 and 11 with radiation diodes and activation foils and the remaining slots with test structures for the HEC power supplies. The setup is shown in Figure 3.

The total collected proton charge in the test was 2.0 C. The neutron flux profile has been measured with sheets of EBT2 radiation film placed for 5 minutes at a reduced machine current of 1.2 $\mu A$ in front of slot 1 and behind all other slots with the exception of slots 7, 10, and 11. During a dedicated fluence test run in 2010 with radiation monitors in each slot the following empirical relation of the 1 MeV equivalent neutron fluence in silicon and the machine current as function of the slot number has been found [5]:

$$\text{fluence} = \frac{1.61 \times 10^{10} n}{\text{cm}^2 \mu \text{C}} \left( \frac{\text{slot} + \text{offset} \pm 2 \text{ mm}}{17 \text{ mm}} \right)^{-2.11},$$  

where slot is the slot-number (1-16), offset is the distance of the first slot from the target (3 mm) and the positions are known with $\pm 2$ mm precision.

Figure 4 shows this fluence curve superimposed by the relative fluence measurements from the radiation films, normalized to the expected fluence in slot 3.

ionization chamber current measurements in 1 s intervals and the conversion factor to the proton beam current was used to calculate the accumulated proton charge as function of time. Together with the beam width and position measurements from the radiation films the fluence for the devices under test in each slot could be derived. The alignment of the setup was done prior to irradiation by means of two lasers indicating the horizontal and vertical planes containing the beam direction. The laser system was however misaligned for the horizontal x and y position of the beam profile along the beam line, the offset of the beam in y, and the expected losses due to nuclear reactions in the devices under test, the flux was found to be the same for all slots.

Figure 2 shows the horizontal and vertical profiles of the proton beam as measured by one of the radiation films. Dips in the data distribution are caused by the soldered legs of housings in the upstream slots. The superimposed Gaussian fits show that the profiles are roughly Gaussian in the area of the ASIC.
The S-parameters have been measured for the frequency range 300 kHz to 100 MHz with a vector network analyzer (VNA) connected via 40 m long cables to the devices under test. Each cable pair has been individually de-embedded with the VNA calibration utility moving the plane of calibration to the input connectors of the boards. The most important parameters are the input port reflection coefficient \( S_{11} \), which can be translated into the input impedance \( Z_{in} = 50 \Omega (1+S_{11})/(1-S_{11}) \), for our case of vanishing feedback coefficient \( S_{12} \) and low load reflection, and the forward transmission coefficient \( S_{21} \). The product of the two gives the transimpedance gain in the frequency domain. This can be interpreted as the input current to output voltage signal amplification. It was found that the effect of protons in the GaAs ASICs is similar to that of neutrons with \( \sim 5 \) times larger fluence. Since the expected ratio of neutron to proton fluence for the HL-LHC is \( \sim 10 \), the neutron tests provide the most limiting numbers.

V. PULSE MEASUREMENTS

In the neutron test triangular voltage pulses similar to the ionization currents in the LAr gaps in ATLAS with 10 different amplitudes between 0.7 mV and 70 mV and pulse lengths of 500 ns have been applied over the 40 m cables at the BB96 ASICs and both the input and the output pulse have been measured with an oscilloscope (the output pulse again after 40 m cable). These measurements were done alternatingly with the S-parameter measurements and provide a more direct way of studying the performance of the electronics. Good agreement with the S-parameter measurements was found for the gain as a function of neutron fluence. The largest expected ionization currents in ATLAS are 250 \( \mu \)A for a single pre-amplifier and about 500 \( \mu \)A for a full chain of several pre-amplifiers and one summing amplifier. Up to these values the linearity of the pre-amplifiers and the full system was measured with the pulses. To convert from input voltage pulses to currents the input impedance and cable damping is taken into account. Thus the voltage range of 0.7 mV to 70 mV corresponds to an input current range of 10 \( \mu \)A to 1000 \( \mu \)A for non-irradiated chips.

VI. WARM IN-SITU RESULTS

The forward transmission coefficient normalized to the value before irradiation and evaluated in the frequency range of the shaper electronics in ATLAS (4 MHz < \( f < 10 \) MHz) as measured in situ for a system of pre-amplifier and summing amplifier during the proton (top) and neutron (bottom) irradiations vs. the fluence. The neutron fluence is given in 1 MeV equivalent for silicon. The red vertical lines indicate HL-LHC limits including a safety factor of \( 5 \). In the neutron case the fluence is again the expected silicon 1 MeV fluence but corrected by the expected differences in GaAs. The colored data points are from the various devices under test and the spread of the points indicate sample variations. Since for the neutron test the fluence varied with slot number the point density and maximal fluence per slot varies.
for protons and up to $\simeq 4 \times 10^{14} \text{n/cm}^2$ for neutrons. The output impedance shows similar behavior and stays roughly twice as long close to 50 $\Omega$. Beyond these limits the impedances stay real but rise quickly resulting in non-matched conditions in ATLAS, where $0.5 - 2 \text{ m} 50 \Omega$ cables connect the detector and the pre-amplifiers and $\simeq 10 \text{ m} 50 \Omega$ cables carry the signals from the ASIC output to the shaper electronics.

VII. COLD RESULTS

About three months after the neutron irradiation tests at NPI in Úžhorod the irradiated boards have been shipped to the lab in Munich and re-tested with the same set of measurements as in situ in both warm and cold (liquid nitrogen) conditions. The three front-most boards received neutron fluences in excess of $2.8 \times 10^{15} \text{n/cm}^2$ (Si 1 MeV eq. in Úžhorod). From Figure 5, it can be seen that beyond $2.1 \times 10^{15} \text{n/cm}^2$ the ASICs approach 0 voltage gain and therefore the three first boards are dead. The board in the fourth slot was exposed to $1.6 \times 10^{15} \text{n/cm}^2$ and this one plus the following four boards from slots 5, 6, 8, and 9, with respective fluences of 10.0, 6.9, 3.8, and $3.0 \times 10^{14} \text{n/cm}^2$ are still operational in warm conditions and still show the same performance as was measured at the end of the run in situ at Úžhorod. In addition to these 5 fluence points non-irradiated ASICs have been measured in warm and cold as well to provide data at zero fluence. Immersed in liquid nitrogen to emulate the conditions inside the ATLAS cryostat the performance is dramatically worse than in warm conditions. Figure 6 shows the normalized forward transmission coefficient and the input impedance for the systems comprised of one pre-amplifier and a summing amplifier as function of neutron fluence for both warm and cold conditions. The sparse data points on the left indicate the cold measurements, while the denser groups of points extending to larger fluences are the measurements in the warm. The performance in the cold is roughly equivalent to that in warm at 3 times larger fluences. The performance degrades quickly beyond $3 - 4 \times 10^{14} \text{n/cm}^2$.

The pulse measurements as described in section VII have also been repeated with the already irradiated boards in warm and in cold. The pulse amplitudes have been adjusted to lower levels to still correspond to $10 - 1000 \mu \text{A}$ input currents for non-irradiated ASICs for the short $(\sim 4 \text{ m})$ cables used at MPP. The increase of non-linearity of the response was found to be described by a power-law relation of input (in) and output (out) amplitude of the form:

$$\text{out} = c \times \text{in}^{x},$$  \hspace{1cm} (2)

where $c$ is a fluence dependent constant and $x$ the power-law exponent. Values for $x$ deviating from 1.0 indicate non-linear behavior.

Figure 7 shows this exponent as function of neutron fluence in warm and cold. Like in the case of the forward transmission coefficient the non-linearities for irradiated ASICs are worse in the cold and correspond roughly to those in the warm at 3 times larger fluences. Since up to 16 pre-amplifiers are summed by one summing amplifier in ATLAS the non-linearities for the pre-amplifiers can not be corrected. An exponent of 1.03 or 0.97 would correspond to a maximum deviation of the output amplitude from linear expectation normalized to the maximum amplitude of 1%. This original
requirement limit for the non-linearity of the pre-amplifiers is indicated by the horizontal lines in Figure 7. Beyond neutron fluences of $3 - 4 \times 10^{14} \text{n/cm}^2$ the non-linearity quickly rises above this level.

The degraded S-parameter measurements in the full frequency range of 300 kHz−100 MHz have been used to predict the response to a typical ionization signal in the HEC after the full electronics chain for the HEC in ATLAS [6]. Figure 8 shows the simulated response to a 400 ns triangular current in a typical HEC channel at 4 different fluence values. It can be seen that on top of the gain degradation the pulse shape worsens with increasing fluence leading to a slower rise- and fall-time, a broader and more delayed positive peak.

VIII. Fluence Conversion for GaAs

The neutron fluence values given in all the figures presented so far are 1 MeV equivalent fluences evaluated for the neutron spectrum at NPI in Rež [4]. In order to compare these values to the HL-LHC the fluences have to be converted to ATLAS spectrum at NPI reported here can thus be converted to ATLAS at the position of the ASICs [4]. Assuming that the results in the cold would be worse compared to the warm in-situ measurements by the same factor of 3 as observed in the neutron tests and thus cold measurements of the proton irradiated ASICs are not available. Assuming that

1) the results in the cold would be worse compared to the warm in-situ measurements by the same factor of 3 as observed in the neutron tests
2) the effect of 1 p/cm² with kinetic energy of 200 MeV in GaAs according to the Kerma factor and specific energy loss [9] can be split in 1.4 n/cm² 1 MeV equivalent neutron fluence in GaAs and 465 pGy total ionization dose (TID)

it is possible to convert the GaAs numbers back to Si again and estimate the impact of the expected Si TID of 6.2 kGy in ATLAS at the HL-LHC, which would be equivalent to $\sim 10\%$ of the expected degradation due to 1 MeV equivalent neutrons.

IX. Conclusions

Irradiation tests of the current readout electronics of the Hadronic Endcap Calorimeter of ATLAS with neutrons and protons with fluence levels exceeding those relevant for the high luminosity upgrade of the LHC have been performed. The most stringent limits have been observed under neutron irradiation and the measurements at liquid nitrogen temperatures show a $\sim 3$ times faster performance degradation compared to room temperature measurements.

The performance limit in the cold was reached at $3 - 4 \times 10^{14} \text{n/cm}^2$ 1 MeV equivalent in Si for neutrons in Rež, where the forward transmission coefficient reduces to $\sim 70\%$, the input (output) impedance increases by $\sim 20\%$ ($\sim 40\%$) above the nominal 50 $\Omega$, and the power-law exponent reaches $\sim 1.03$. These levels correspond to $4 - 5.4 \times 10^{14} \text{n/cm}^2$ 1 MeV equivalent in Si for ATLAS at the position of the ASICs.

With the current safety factor of 5 for the radiation level simulations for ATLAS this limit will be reached at the HL-LHC. The impact on the ATLAS physics performance of the degraded electronics and upgrade options of the readout electronics for the HL-LHC are currently under study.

References

[1] D. M. Gingrich et al., Construction, assembly and testing of the ATLAS hadronic end-cap calorimeter, JINST 2 P05005 (2007).
[2] The ATLAS Collaboration, The ATLAS Experiment at the CERN Large Hadron Collider, JINST 3 S08003 (2008).
[3] J. Ban et al., Cold Electronics for the Liquid Argon Hadronic End-cap Calorimeter of ATLAS, Nucl. Instr. & Meth. A556, 158-168 (2006).
[4] P. Bém et al., The NPI cyclotron-based fast neutron facility, International Conference on Nuclear Data for Science and Technology 2007, 555-558, http://dx.doi.org/10.1051/ndata:07598 (2007).
[5] D. Dannheim, n-flux measurement Rež irradiation April 2010, http://indico.mppmu.mpg.de/indico/getFile.py/access?contribId=3&resId=1&materialId=slides&confId=834 (2010).
[6] L. Kurchaninov, Modeling of the HEC Electronics Chain, ATLAS HEC-Note-109 (2001).
[7] A. Vasilescu and G. Lindstroem, Displacement damage in silicon, online compilation, http://sesam.desy.de/members/gunnar-Si-efflux.html.
[8] P. J. Griffin et al., SAND92-0094 (1993), A. Konobeyev, J. Nucl. Mater. 186-117 (1992), G.P. Summers et al., IEEE NS 40, 1372 (1993), M. Huh and P.A. Aarnio, Nucl. Instr. & Meth. A335, 580 (1993), A.M. Ougoug et al., IEEE NS 31, 2219 (1990), A.L. Barry et al., IEEE NS 42, 2104 (1995).
[9] S. Baranov et al., Estimation of Reaction Background, Impact on Detectors, Activation and Shielding Optimization in ATLAS, ATL-COM-2005-001 (2005).
[10] J. Beringer et al. (Particle Data Group), 2012 Review of Particle Physics, Phys. Rev. D86, 010001 (2012).