Universal ion-transport descriptors and classes of inorganic solid-state electrolytes

Cibrán López,1,2,3 Agustí Emperador,1 Edgardo Saucedo,2,4 Riccardo Rurali,3 and Claudio Cazorla1,2

1Departament de Física, Universitat Politècnica de Catalunya, 08034 Barcelona, Spain
2Barcelona Research Center in Multiscale Science and Engineering, Universitat Politècnica de Catalunya, 08019 Barcelona, Spain
3Institut de Ciència de Materials de Barcelona, ICMAB-CSIC, Campus UAB, 08193 Bellaterra, Spain
4Department of Electronic Engineering, Universitat Politècnica de Catalunya, 08034 Barcelona, Spain

Solid-state electrolytes (SSE) with high ion conductivity are pivotal for the development and large-scale adoption of green-energy conversion and storage technologies such as fuel cells, electrocatalysts and solid-state batteries. Yet, SSE are extremely complex materials for which general rational design principles remain indeterminate. Here, we unite first-principles materials modelling, computational power and modern data analysis techniques to advance towards the solution of such a fundamental and technologically pressing problem. Our data-driven survey reveals that the correlations between ion diffusivity and other materials descriptors in general are monotonic, although not necessarily linear, and largest when the latter are of vibrational nature and explicitly incorporate anharmonic effects. Surprisingly, principal component and k-means clustering analysis show that elastic and vibrational descriptors, rather than the usual ones related to chemical composition and ion mobility, are best suited for reducing the high complexity of SSE and classifying them into universal classes. Our findings highlight the need of considering databases that incorporate temperature effects to improve our understanding of SSE and point towards a generalized approach to the design of energy materials.

Social networks use modern data analysis techniques to improve their customers experience and increase advertising revenues [1]. Each mouse click and fingers action on the touchscreen reveal information on the users preferences that can be employed to classify individuals into similarity groups and thus better select the contents they are exposed to. Materials, in analogy to humans, conform to highly diverse and complex collectives and as such advanced data analysis techniques are being increasingly applied on them to improve their design and recommend possible uses [2, 3]. A necessary condition for the meaningful development and application of data-driven materials design strategies is the existence of comprehensive and reliable databases.

Solid-state electrolytes (SSE) are a class of energy materials in which specific groups of ions may start to diffuse throughout the crystalline matrix driven by the thermal excitations [4]. SSE are the pillars of green-energy conversion and storage technologies like fuel cells, electrocatalysts and solid-state batteries, hence tuning of their ion-transport properties turns out to be critical for the fields of Energy and Sustainability. SSE, however, are highly complex materials that present disparate compositions, structures, thermal behaviors and ion mobilities, thus it is difficult to ascribe them to general and rational design principles. These difficulties have motivated researchers to seek for easy to measure (or calculate) quantities that may serve as good descriptors of the ion conductivity; examples of such descriptors include structural parameters, defect formation energies, atomic polarizabilities and lattice dynamics [5, 6]. In recent years, pinpointing the role of phonon dynamics on ion transport has attracted special and increasing attention. Actually, for some specific SSE it has been demonstrated that lattice anharmonicity is one of the most influential factors affecting their ion mobility [7, 8].

Quantum mechanics-based density functional theory (DFT) [9] has proven tremendously successful in the field of computational materials science, and currently several databases of automated DFT calculations are being widely employed for materials design applications [10–19]. Nevertheless, despite of their great successes, the existing DFT databases might not be entirely adequate for progressing in the design and understanding of SSE because they mostly contain information generated at zero temperature (e.g., structural parameters and formation energies) and thus completely disregard anharmonicity and T-induced effects [20]. In addition, modern high-throughput and machine learning studies relying on such DFT databases mainly have targeted Li and Na-based SSE families due to their predominance in electrochemical storage applications [8, 21, 22]. To holistically better understand the phenomena of ion transport, however, it might be necessary to analyse in equal measure other classes of SSE, like those involving mobile O, Cu, Ag and halide ions, which are technologically relevant as well [23, 24].

Here, we present a data-driven analysis of SSE that covers aspects generally unaddressed by previous computational studies and the existing DFT materials databases. First, a comprehensive first-principles database was created for prototypical families of inorganic SSE containing both sets of zero-temperature DFT and finite-temperature ab initio molecular dynamics (AIMD) results. Subsequently, a thorough correlation study between the ion diffusion coefficient (D) and other materials features was performed to determine universal
ion-transport descriptors (as well as those specific to Li-based SSE). By relying on this new knowledge and the introduced DFT-AIMD database, several machine learning models were trained for the prediction of $D$ and other $T$-dependent quantities. Finally, principal component and k-means clustering analysis, data techniques customarily employed in the social sciences, were applied to reduce the high complexity of the SSE landscape and determine universal classes of fast-ion conductors.

**Curated first-principles SSE database.** The generated SSE DFT-AIMD database [26] comprises a total of 61 materials of which 46% contain Li, 23% halide (i.e., F, Cl, Br and I), 15% Na, 8% O and 8% Ag/Cu atoms as the mobile ions. These percentages were selected in order to roughly reproduce the relative abundances of fast-ion conductors reported in the literature [27]. The generated SSE DFT-AIMD database contains materials with both
Stoichiometric and non-stoichiometric compositions and the AIMD results were obtained over a broad range of temperatures (Supplementary Tables 1–3 and [20]).

To analyze the degree of similarity between all the surveyed SSE, a great variety of descriptors were estimated for each material adding up to a total of 54 (the complete list of descriptors is detailed in the Methods section). Some of these descriptors had been already proposed in the literature (e.g., band gap and vacancy formation energy) while some others were totally new (e.g., harmonic phonon energy and Pugh’s modulus ratio). The descriptors were classified into three general categories: “mechanical-elastic”, “diffusive-vibrational” and “structural-compositional”. The value of some descriptors were obtained from zero-temperature DFT calculations (“mechanical-elastic” and “structural-compositional”) while the rest (“diffusive-vibrational”) were deduced from AIMD simulations performed at temperatures above ambient (Methods and Supplementary Tables 1–3).

It is worth noting that the results obtained from AIMD simulations explicitly account for anharmonic effects, which constitutes one of the most important novelties and technical advances of the present work and introduced SSE database. Moreover, most vibrational descriptors were estimated considering the following cases (1) all the ions, (2) only non-diffusive ions and (3) only diffusive ions, in order to better substantiate the role of the vibrating crystalline matrix on ion transport (Methods). The approximate computational cost of the generated SSE DFT-AIMD database was of 50 Million CPU hours.

Correlations between pairs of SSE descriptors. The correlation for a couple of materials descriptors, x and y, can be quantified in several non-unique ways [25]. In this work, we considered the Pearson ($c_p$) and Spearman ($c_S$) correlation coefficients which are defined like:

\[
c_p(x, y) = \frac{\text{cov}(x, y)}{\sigma_x \sigma_y} \quad \text{and} \quad c_S(x, y) = c_p \left[ R(x), R(y) \right],
\]

where $\sigma_i$ is the standard deviation of the descriptor i and $R(i)$ the rank of the i samples. The covariance function is expressed as:

\[
\text{cov}(x, y) = \langle xy \rangle - \langle x \rangle \langle y \rangle,
\]

where $\langle \cdot \rangle$ denotes expected value. The Spearman correlation coefficient is able to detect monotonic dependencies between pairs of descriptors while the Pearson can only identify linear correlations. Thus, the $c_S$ correlation coefficients are more general and robust than $c_p$ (i.e., can assess monotonic relationships whether linear or not). For this important reason, and despite the fact that linear correlations have been assumed in most previous SSE studies [7, 9], we will stick to the Spearman correlation definition for the rest of our analysis.

Figure 1 shows the Spearman correlation coefficients estimated for all pairs of materials descriptors considering the entire DFT-AIMD database (an analogous Pearson correlogram can be found in the Supplementary Fig.1). In view of the preeminence of Li-based SSE in electrochemical applications, the same correlation analysis was performed for this family of materials alone (Fig. 2(d)). To assess the statistical significance of the estimated $c_S$ correlograms, we computed the corresponding $p$-value matrices (Figs. 1b, d). The $p$-value represents the probability for a particular correlation result to arise if the null hypothesis (i.e., no correlation at all) were true, thus the smaller the calculated $p$-value the more statistically significant $c_S$ is.

In a bird’s eye view, the two correlograms obtained for all SSE and only those containing Li ions look quite similar. Nevertheless, the $p$-value matrix estimated for all SSE displays a noticeably higher number of statistically significant cases (arbitrarily defined here as $p < 0.2$), probably due to the larger amount of samples. Reassuringly, a number of already expected high correlation coefficients, like those estimated for couples of vibrational and elastic quantities that are physically related (e.g., $F_{\text{vib}}$ and $S_{\text{vib}}$), emerge from the calculated $c_S$ maps. For the sake of focus, hereafter we will concentrate on the correlations involving the ion diffusion coefficient ($D$).

Figure 2 encloses a standardized representation [that is, $\tilde{x} \equiv (x - \langle x \rangle) / \sigma_x$] of the pairs of descriptors $D-C_v$ and $D-\langle \omega \rangle$, where $C_v$ stands for the lattice heat capacity and $\langle \omega \rangle$ the average vibrational frequency (Methods). In these two cases, as well as in others not shown here, it is clearly appreciated that the dependency between $D$ and the other quantities is far from linear although roughly monotonic. This outcome confirms that for determining reliable relationships between SSE features the Spearman correlation analysis is certainly more suitable than the usual Pearson approach. Actually, there are significant discrepancies between the Spearman and Pearson correlation maps; for instance, $c_S$ amounts to $-39\%$ for the pair of descriptors $D-\langle \omega \rangle$ (Fig. 1k) whereas $c_p$ renders a significantly smaller value of $-23\%$ (Supplementary Fig. 2).

Universal ion diffusion descriptors. Figure 3 shows the Spearman correlation coefficients estimated for all pairs of descriptors involving $D$ and considering the entire DFT-AIMD database. All the AIMD-based vibrational and diffusive descriptors were estimated at $T = 500 \pm 100$ K. First, we note that larger $|c_S|$ values are associated with statistically more significant correlation results (i.e., smaller $p$-values). And second, the estimated correlation coefficients in general are not very high: only 19 out of the 53 pairs of materials descriptors present $|c_S|$’s larger than 20% while the maximum correlation value only amounts to 39% (obviously, the $D-D$ pair was excluded here). These low-correlation outcomes are consistent with the usual difficulties encountered in the settlement of flawless ion transport descriptors [9].
FIG. 2. Correlation study of the ion diffusion coefficient with other materials descriptors. (a) Standardized representation of the ion diffusion coefficient $D$ along with other materials descriptors. The descriptors correlations are, to some extent, monotonic but not linear as it is shown by the orange and blue lines therein (both simple guides to the eyes). Spearman correlation coefficients for $D$ and the rest of materials descriptors considered in this study, obtained by taking into account (b) all and (c) exclusively the Li-based compounds included in our DFT-AIMD database. The $p$-value results corresponding to the Spearman correlation coefficients are indicated with different colours. All the AIMD-based diffusive and vibrational descriptors were estimated at $T = 500 \pm 100$ K.
Interestingly, the largest $D$ correlations are found for AIMD-based vibrational descriptors (Methods) like the phonon band center (or average lattice frequency), $\langle \omega \rangle$ (−39%), lattice heat capacity, $C_V$ (+39%), vibrational free energy, $F_{\text{vib}}$ (−37%), and vibrational entropy, $S_{\text{vib}}$ (+33%). These results indicate that insulator materials with small average phonon frequencies, large heat capacities and large vibrational entropies should be good ion conductors. It is worth noticing that strongly anharmonic materials perfectly fit into this description, thus our data-driven results generalize the conclusions of recent experimental SSE studies revealing that low-energy phonon modes can actively influence ion diffusion in some specific materials [9–14].

Our correlation analysis provides further valuable insights. First, when the vibrational descriptors were estimated considering either non-diffusive or diffusive ions alone (superscripts “nd” and “d” in Fig. 2, respectively) the value of the $D$ correlation coefficients slightly decreased in the first case ($|c_S| = 30\%$) and practically vanished in the second (except that corresponding to $\langle \omega_3 \rangle$). This outcome highlights the existence of a strong and general interplay between the vibrating crystalline matrix and mobile ions. And second, when considering vibrational descriptors that do not explicitly take into account anharmonic effects, like the lowest-energy optical phonon mode calculated at $T = 0$ K (Γ in Fig. 2), the resulting $D$ correlation coefficient (−11%) significantly drops in comparison to those obtained for anharmonic quantities (besides, the corresponding $p$-value increases). Thus, scrutiny of anharmonicity appears to be indispensable for the evaluation of reliable and statistically meaningful $D$ correlation coefficients.

Few descriptors belonging to the “structural-compositional” category also correlate appreciably high with $D$. Of special mention are the vacancy formation energy of the mobile ions ($E_{\text{vac}}$, −22%), the crystal polarizability ($\alpha_C$, +25% -calculated with the Clausius-Mossotti relation–) and the symmetry of the perfect lattice (SO, +27%) [29]. On the other hand, intrinsically electronic properties like the energy band gap ($E_g$) and dielectric constant ($\varepsilon$) have virtually no correlation with the ion diffusivity ($|c_S| \leq 5\%$). As a word of caution, we note that when the correlations between $D$ and other materials descriptors are assumed to be linear (i.e., Pearson’s approach) the resulting conclusions significantly differ from those just explained (Supplementary Fig. 2). In particular, most $D$ correlation coefficients turn out to be smaller than the corresponding Spearman values and the materials descriptors belonging to the “mechanical-elastic” category (e.g., the Young and shear moduli – $E$ and $G$ –) become equally relevant than the vibrational features.

Figure 2 shows the Spearman $D$ correlation coefficients estimated exclusively for Li-based SSE. Intriguingly, the resulting $c_S$ chart differs appreciably from that estimated considering the entire DFT-AIMD database (Fig. 3). First, the $D$ correlation coefficients in general present larger values with a total of 11 pairs of materials descriptors scoring above 40%. Some of the largest $|c_S|$ correspond to the AIMD-based vibrational descriptors $F_{\text{vib}}$ (−42%), $S_{\text{vib}}$ (+42%) and $\langle \omega_3 \rangle$ (−63%). However, in contrast to the all-SSE case, now $\Gamma$, which is estimated at $T = 0$ K and does not explicitly account for anharmonicity, is strongly correlated with $D$ as well (−47%). Moreover, several descriptors belonging to the “mechanical-elastic” category that, to the best of our knowledge, have not been previously proposed in the literature like the Vickers’ hardness, $H_V$ (−43%), Pugh’s modulus ratio, $\kappa$ (−56%), Poisson’s ratio, $\nu$ (+55%), Cauchy’s pressure, $P_C$ (+48%), and velocity ratio, $v_v$ (+56%), now also render very high $|c_S|$ values. Therefore, in terms of key $D$ descriptors, Li-based compounds are plainly different from the average SSE, a finding that fundamentally justifies the large number of studies focusing on the ion transport properties of this family of materials.

**Machine learning models for prediction of $T$-dependent properties.** In view of the complex relationships between $D$ and other materials descriptors (Fig. 2), several machine learning (ML) models based on artificial neural networks were trained in our SSE DFT-AIMD database with the aim of predicting the ion diffusion coefficient and other relevant $T$-dependent properties of SSE such as $\langle \omega \rangle$ and $C_V$ (Methods). We considered two different ML training schemes: (1) considering all the materials descriptors (denoted as “anharmonic”) and (2) excluding the AIMD-based vibrational descriptors (“harmonic”). The predictions of our trained ML models for a validation set of 12 compounds are shown in Fig. 3. Therein, it is appreciated that the two trained ML models can predict the finite-temperature values of $\langle \omega \rangle$ and $C_V$ with high accuracy. In particular, the mean absolute percentage error (MAPE) of the “anharmonic” (“harmonic”) ML model amounts to 2.5% (7.5%) and only 0.5% (1.9%) for $\langle \omega \rangle$ and $C_V$, respectively. In stark contrast, the ML predictions for the ion diffusion coefficient are much less accurate and there is a huge difference in the level of precision achieved with the “anharmonic” (MAPE of 60%) and “harmonic” (290%) ML models.

Several conclusions follow from the ML results enclosed in Fig. 3. First, the SSE DFT-AIMD database introduced in this work appears to be comprehensive enough to ensure proper training of ML models able to make accurate predictions of certain $T$-dependent materials properties. And second, ML-based prediction of the ion diffusivity appears to be a particularly challenging task. In this latter case, however, a big improvement is achieved when AIMD-based anharmonic vibrational descriptors are explicitly incorporated into the ML model (also in the $\langle \omega \rangle$ and $C_V$ cases). This outcome indirectly corroborates our previous finding that anharmonicity is a key general factor influencing ion transport. Nonetheless, to improve the “anharmonic” ML prediction of $D$ probably it is necessary to increase
FIG. 3. Machine learning (ML) models trained in our DFT-AIMD database for prediction of different SSE $T$-dependent quantities. The ML models were trained by considering and neglecting AIMD-based vibrational descriptors that explicitly incorporate anharmonic effects, labelled as “anharmonic” and “harmonic”, respectively. (a) First momentum of the vibrational density of states obtained from AIMD simulations, $\langle \omega \rangle$. (b) Constant volume heat capacity obtained from AIMD simulations, $C_V$. (c) Ionic diffusion coefficient obtained from AIMD simulations, $D$. “MAPE” stands for the mean absolute percentage error of the ML predictions.
FIG. 4. Principal component analysis results obtained for the SSE DFT-AIMD database. (a) Eigenvalues corresponding to the diagonalization of the Spearman correlation matrix obtained by considering the entire DFT-AIMD database. (b) Eigenvector components of the first three principal components obtained from the diagonalization of the Spearman correlation matrix obtained by considering the entire DFT-AIMD database.

The number of SSE materials and descriptors in our DFT-AIMD database and/or resort to alternative and more advanced ML approaches (e.g., graph neural networks [30]).

Complexity reduction in the SSE landscape. Principal component analysis (PCA) is a statistical technique widely employed for analyzing large datasets containing a high number of features. PCA increases the interpretability of a dataset by reducing its dimensionality and simultaneously preserving the maximum amount of information. Complexity reduction is accomplished by linearly transforming the data into a new coordinate system where most of its variation can be described with fewer dimensions. The principal components are the eigenvectors of the dataset correlation matrix, which are expressed as linear combinations of the initial descriptors. The first principal component, the one with the largest eigenvalue, maximizes the variance of the projected data. The $i$-th principal component corresponds to a direction that is orthogonal to the previous $i-1$ principal components and along which the variance of the projected data is maximized as well.

Figure 4 shows the results of diagonalizing the Spearman correlation matrix obtained for the entire SSE DFT-AIMD database. The first three principal components (PC) account for about two thirds of the total variance in the original 54-dimensional dataset (as quantified by the sum of their normalized eigenvalues, $\approx 62\%$) hence its complexity can be greatly reduced by considering data projections on the orthogonal three-dimensional space PC1–PC2–PC3. PC1 presents a dominant “mechanical-elastic” character, PC2 “vibrational” and PC3 “structural” (Fig 4b). Intriguingly, the contribution of the ion diffusivity to each of these PC’s is practically zero, namely, 0.2% to PC1, 0.8% to PC2 and
FIG. 5. K-means clustering analysis results obtained for the SSE DFT-AIMD database. (a) Classification of the analyzed materials in the orthogonal bidimensional space PC1–PC2. (b) Materials population of each group identified in the PC1–PC2 space expressed in terms of the mobile ion species. (c) Classification of the analyzed materials in the orthogonal tridimensional space PC1–PC2–PC3. (d) Materials population of each group identified in the PC1–PC2–PC3 space expressed in terms of the mobile ion species. To improve visual clarity, some points have been removed from the plots without affecting the main conclusions.

1.3% to PC3. This data-driven outcome indicates that when it comes to characterize the great disparity of SSE, with the aim of fundamentally better understanding them and to establish general SSE categories, the ubiquitous $D$ descriptor is actually irrelevant. Likewise, the compound stoichiometry ($Stc$) and dielectric constant ($\epsilon$) hardly contribute to the first three PC’s hence they neither can be regarded as universally distinctive SSE features. By contrast, elastic and vibrational descriptors like $E$, $H_V$, $\langle \omega \rangle$ and $C_V$ become most pertinent for the evaluation of SSE similarities and general classification purposes.

K-means clustering analysis. Figure 5 encloses the results of our k-means clustering analysis performed for the entire SSE DFT-AIMD database. K-means clustering is an unsupervised learning algorithm that classify sets of objects in such a way that objects within the same group, called “cluster”, are more similar to each other in a broad sense than to the objects in other clusters. We selected a subminimal number of 7 clusters to account for the SSE database variance based on the outcomes of the elbow and silhouette methods (Supplementary Figs. 3–4). This number of clusters is already larger than the number of $A$-based SSE families considered in this study (i.e., 6 with $A = Li, Na, halide, Ag, Cu$ and $O$). Thus, it straightforwardly follows that the materials composition, despite of its obvious utility in naming compounds, should not be regarded as a fine descriptor of SSE diversity since, at least, one SSE family will spread over more than one k-means cluster.

Figures 5a–b show the results of our k-means clustering analysis performed in the simplified PC1–PC2 space. It is noted that Li-based SSE are present in 5 out of...
the total 7 clusters. From those 5 clusters, Li-based SSE are the most abundant in 80% of the cases and overall they share similarities with other Na, halide and O-based SSE (although not necessarily in terms of ion conductivity). In clusters number 1 and 2, which are respectively characterized by dominant PC1 (“elastic”) and PC2 (“vibrational”) components, Li-based SSE actually conform the entire population. From these outcomes, we may readily conclude that (1) Li-based SSE are intrinsically different from Ag- and Cu-based SSE, which in turn are highly similar because inhabit the same cluster, and (2) Li-based SSE can be partitioned into several similarity subgroups attending to their elastic and vibrational properties. Likewise, halide-based SSE appear in 4 different clusters, Na-based in 3 and O-based in 2. Thus, as it was mentioned above, chemical composition is not a good descriptor for grouping SSE into similarity categories.

Figures 5c–d enclose the k-means clustering results obtained in the expanded PC1–PC2–PC3 space. In this case, the main findings are very similar to those just explained for the reduced P1–P2 space, namely, Li-based SSE are present in 5 out of the total 7 clusters and they are particularly numerous in the majority of those groups (e.g., 100% in cluster number 7 and 67% in cluster number 2). Likewise, halide-based SSE spread over 4 different clusters, Na-based over 4, O-based over 2 and Cu/Ag-based only appear in 1. The Li-based SSE family overall shares similarities with other Na, halide and O-based SSE (not so with Cu- and Ag-based SSE), and most subgroup differences (i.e., relative distances between clusters centroids) are contained in the P1–P2 plane. Thus, the PC3 (“structural”) dimension does not appear to add sensible information on SSE diversity and for grouping purposes is practically expendable (in accordance with its relatively small eigenvalue of ≈ 4%, Fig 5c).

The presented k-means clustering analysis enlightens the difficulties encountered in the rational design of SSE with specific ion mobility. The bulk of the variation in the SSE family is encoded in the materials elastic and vibrational properties, neither in the ion mobility nor their chemical composition. This finding implies that materials which can be rigorously considered as overall highly similar (because they belong to a same k-means cluster) in practice may exhibit very different ion diffusion and chemical features (e.g., Li-based and halide-based SSE). Conversely, materials which render very similar ion mobilities and chemical compositions (e.g., Li-based SSE inhabiting groups 7 and 3 in Fig 5c) may behave radically different in terms of other measurable quantities. These conclusions are consistent with the D correlation results enclosed in Fig 2 which show that Li-based SSE can significantly depart from the general trends averaged over all SSE.

In summary, we have presented an original and comprehensive SSE data-driven study on the correlations of the ion diffusion with other materials descriptors as well as a rigorous examination of universal SSE categories, based on a new and thorough DFT-AIMD database comprising both zero-temperature and finite-T first-principles results. It has been demonstrated that ion diffusion correlates strongly and monotonically, not necessarily linearly, with vibrational descriptors that explicitly incorporate anharmonic effects (i.e., are estimated from AIMD simulations). In the particular case of Li-based SSE, the ion mobility also correlates significantly with elastic quantities like the Vickers’ hardness, Pugh’s modulus ratio, Poisson’s ratio and Cauchy’s pressure, pertinent ion-diffusion descriptors that previously have been overlooked in the literature. Furthermore, most of the variation in the generated SSE 54-fold dimensional space can be resolved in terms of elastic and vibrational descriptors; ion mobility and chemical composition are very much irrelevant when it comes to quantify the SSE diversity, a fact that complicates the rational design of SSE with targeted ion conductivities. The present data-driven study highlights the necessity to consider finite-temperature effects in a high-throughput fashion to better understand SSE and improve the predictions of machine learning models in them; it also provides new theoretical guidelines for analyzing materials that in analogy to SSE are highly anharmonic and technologically relevant (e.g., thermoelectrics and superconductors).

**METHODS**

**First-principles calculations outline.** *Ab initio* calculations based on density functional theory (DFT) were performed to analyse the physico-chemical properties of bulk SSE. We performed these calculations with the VASP code [31] by following the generalized gradient approximation to the exchange-correlation energy due to Perdew et al. [32]. (For some halide compounds, possible dispersion interactions were captured with the D3 correction scheme developed by Grimme and co-workers [33].) The projector augmented-wave method was used to represent the ionic cores [34] and for each element the maximum possible number of valence electronic states was considered. Wave functions were represented in a plane-wave basis typically truncated at 750 eV. By using these parameters and dense k-point grids for Brillouin zone integration, the resulting zero-temperature energies were converged to within 1 meV per formula unit. In the geometry relaxations, a tolerance of 0.005 eVÅ−1 was imposed in the atomic forces.

**First-principles molecular dynamics simulations.** *Ab initio* molecular dynamics (AIMD) simulations based on DFT were performed in the canonical (N, V, T) ensemble (i.e., constant number of particles, volume, and temperature) for all the considered bulk materials. The selected volumes were those determined at zero temperature hence thermal expansion effects were
neglected; nevertheless, based on previously reported molecular dynamics tests \cite{12}, thermal expansion effects are not expected to affect significantly the estimation of the ion-transport properties of SSE at moderate temperatures (i.e., $T = 500 \pm 100 \text{ K}$). The concentration of ion vacancies in the non-stoichiometric compounds was also considered independent of the temperature and equal to $\sim 1\text{--}2\%$. The temperature in the AIMD simulations was kept fluctuating around a set-point value by using Nose-Hoover thermostats. Large simulation boxes containing $N_{\text{ion}} \sim 200\text{--}300$ atoms were employed in all the cases and periodic boundary conditions were applied along the three Cartesian directions. Newton’s equations of motion were integrated by using the customary Verlet’s algorithm and a time-step length $\delta t$. The finite-temperature simulations within the first Brillouin zone was employed in all the cases and periodic boundary conditions were applied along the three Cartesian directions. Newton’s equations of motion were integrated by using the customary Verlet’s algorithm and a time-step length $\delta t$. The finite-temperature simulations typically comprised long simulation times of $t_{\text{total}} \sim 100\text{--}200 \text{ ps}$. For each material, we ran an average of 3 AIMD simulations at different temperatures and considering both stoichiometric and non-stoichiometric compositions (Supplementary Tables 1–3 and \cite{26}). Previous tests performed on the numerical bias stemming from the finite size of the simulation cell and duration of the molecular dynamics runs reported in work \cite{12} indicate that the adopted $N_{\text{ion}}$ and $t_{\text{total}}$ values should provide reasonably well converged results for the ion diffusivity and vibrational density of states of SSE.

**Estimation of key diffusive and vibrational properties.** The mean-squared displacement (MSD) was estimated like:

$$\text{MSD}(\tau) = \frac{1}{N_{\text{ion}} (N_{\text{step}} - n_{\tau})} \times \sum_{i=1}^{N_{\text{ion}}} \sum_{j=1}^{N_{\text{step}} - n_{\tau}} |r_i(t_j + \tau) - r_i(t_j)|^2,$$

where $r_i(t_j)$ is the position of the migrating ion $i$ at time $t_j (= j \cdot \delta t)$, $\tau$ represents a lag time, $n_{\tau} = \tau/\delta t$, $N_{\text{ion}}$ is the total number of mobile ions, and $N_{\text{step}}$ the total number of time steps. The maximum $n_{\tau}$ was chosen equal to $N_{\text{step}}/2$ in order to accumulate enough statistics to reduce significantly the fluctuations in MSD(\tau) at large $\tau$’s. The diffusion coefficient then was obtained by using the Einstein relation:

$$D = \lim_{\tau \to \infty} \frac{\text{MSD}(\tau)}{6\tau}.$$

In practice, we performed linear fits over the averaged MSD(\tau) values calculated within the lag time interval $\tau_{\text{max}}/2 \leq \tau \leq \tau_{\text{max}}$.

To estimate the vibrational density of states (VDOS) of bulk SSE considering anharmonic effects, $g(\omega)$, we calculated the Fourier transform of the corresponding velocity-velocity autocorrelation function as obtained directly from the AIMD simulations, namely:

$$g(\omega) = \frac{1}{N_{\text{ion}}} \sum_{i} \int_{0}^{\infty} (\mathbf{v}_i(\tau) \cdot \mathbf{v}_i(0)) e^{i\omega \tau} d\tau,$$

where $\mathbf{v}_i(t)$ represents the velocity of the $i^{th}$ atom at time $t$, and $\langle \cdots \rangle$ denotes statistical average in the $(N,V,T)$ ensemble. Once the density of vibrational states was determined, it was straightforward to calculate the corresponding phonon band center (or average lattice frequency), $\langle \omega \rangle$, defined like:

$$\langle \omega \rangle = \frac{\int_{0}^{\infty} \omega g(\omega) d\omega}{\int_{0}^{\infty} g(\omega) d\omega},$$

which also depends on $T$. Likewise, the contribution of a particular group of ions to the full VDOS was estimated by considering those ions alone in the summation appearing in Eq.(5). In order to determine a characteristic low-energy phonon frequency for bulk SSE, we defined the quantity:

$$\langle \omega_{\text{30}} \rangle = \frac{\int_{0}^{\omega_{\text{max}}} \omega g(\omega) d\omega}{\int_{0}^{\omega_{\text{max}}} g(\omega) d\omega},$$

for which we imposed an arbitrary cut-off frequency of $\omega_{\text{max}} = 30 \text{ meV}$. The analytical expression for other vibrational descriptors (e.g., $E_{\text{vib}}$, $E_{\text{vib}}$ and $C_V$) can be found in work \cite{35}.

**Machine learning models.** The Scikit-learn package in Python \cite{36} was used to encode the non-numeric descriptors as well as to implement the Artificial Neural Network (ANN) conforming our machine learning model. For the generation of the input data, the simulations involving all compounds, compositions and temperatures in our SSE DFT-AIMD database were taken into consideration (i.e., a total of 174 samples, Supplementary Tables 1–3 and \cite{26}). The non-numeric descriptors (i.e., the diffusive chemical element, stoichiometrically, chemical composition of the compound and symmetry of the relaxed structure) were encoded with the one-hot encoding approach, and all input data was normalized using a standard scaler. Specifically, a Multi-Layer Perceptron Regressor (MLPR) was implemented, consisting on input, hidden and output layers. As output layer, the algorithm was defined in such a way that any of the considered descriptors could be used as dependent variable. Consequently, the input layer was constructed as the set of all the other descriptors. Optionally, anharmonic descriptors could be removed from the input layer if desired. Finally, 6 hidden layers of 150, 500, 50, 150, 70 and 100 neurons, respectively, showed the best performance.

Attending to the extraction of metrics, K-fold validation was implemented: on each interation, the model was required to predict the output for one element using the rest as training set. Therefore, given that each element consists of a different number of simulations
(the original dataset presents a variable number of simulated temperatures and stoichiometricities for each element), the computed metrics were weighted with the number of predicted outputs and then divided by the total amount of simulations. The optimization of the model was monitored by using the mean absolute percentage error (MAPE) defined like:

$$ MAPE = \frac{1}{N} \sum_{i=1}^{N} \left| \frac{x_i - \hat{x}_i}{\hat{x}_i} \right|, $$

where $N$ is the total number of samples in the set, $\{x\}$ the predicted outputs and $\{\hat{x}\}$ the actual values in the DFT-AIMD database. Note that these metrics can be extracted from both the training and test sets. As optimal hyperparameters, Adam optimizer with the square error as loss function and constant learning rate of 0.001, rectified linear unit (ReLU) activation function, and $\alpha = 0.05$ strength for the $L^2$ regularization term of the loss function were used.

**SSE descriptors abbreviations.** To analyze the similarities and dissimilarities between fast-ion conductors a great variety of different physical descriptors were estimated for each SSE, which are summarized in Table I. The descriptors are generally classified according to the quality they refer to, in particular: “mechanical-elastic” (M-E), “diffusive-vibrational” (D-V) and “structural-compositional” (S-C). It may be noted that most D-V descriptors like the mean phonon frequency (both with and without cut-off), harmonic phonon energy, constant-volume heat capacity, Helmholtz free energy and entropy, were calculated for the materials as a whole (i.e., considering both diffusive and non-diffusive ions) and also exclusively considering either the non-diffusive (denoted as “nd” in the figures) or diffusive atoms (denoted as “d” in the figures). The total number of descriptors considered in this work is equal to 54. The descriptors estimated from AIMD (DFT) simulations were obtained at $T = 500 \pm 100$ K ($T = 0$ K).

**DATA AVAILABILITY**

The data that support the findings of this study are available upon reasonable request from the authors C.L. and C.C. and the URL: https://superionic.upc.edu/
| Symbol | Descriptor (M-E) | Estimation approach |
|--------|------------------|---------------------|
| $\lambda$ | 1st Lamé parameter | DFT |
| $B$ | Bulk modulus | DFT |
| $E$ | Young modulus | DFT |
| $G$ | Shear modulus | DFT |
| $\nu$ | Poisson’s ratio | DFT |
| $\sigma$ | P-wave modulus | DFT |
| $H_V$ | Vickers’ hardness | DFT |
| $\kappa$ | Pugh’s modulus ratio | DFT |
| $P_C$ | Cauchy’s pressure | DFT |
| $v_l$ | Longitudinal wave velocity | DFT |
| $v_t$ | Transverse wave velocity | DFT |
| $v_r$ | Velocity ratio | DFT |
| $\langle v \rangle$ | Average wave velocity | DFT |

| Symbol | Descriptor (D-V) | Estimation approach |
|--------|------------------|---------------------|
| $\gamma$ | Lindemann ratio | AIMD |
| $\Gamma$ | Lowest-energy optical phonon mode | DFT |
| $\langle \omega \rangle$ | Mean frequency | AIMD |
| $\langle \omega_{30} \rangle$ | Mean frequency (cut-off at 30 meV) | AIMD |
| $E_{vib}$ | Vibrational phonon energy | AIMD |
| $C_v$ | Constant volume heat capacity | AIMD |
| $\theta_D$ | Debye temperature | AIMD |
| $F_{vib}$ | Vibrational Helmholtz free energy | AIMD |
| $S_{vib}$ | Vibrational entropy | AIMD |
| $D$ | Diffusion coefficient | AIMD |
| $msd$ | Mean-squared displacement | AIMD |

| Symbol | Descriptor (S-C) | Estimation approach |
|--------|------------------|---------------------|
| $Z_N$ | Nominal charge | Formula |
| $Z_B$ | Born effective charge | DFT |
| $\epsilon$ | Ion-clamped macroscopic dielectric constant | DFT |
| $M$ | Mobile ion atomic mass | Formula |
| $\alpha_I$ | Mobile ion polarizability | DFT |
| $\alpha_C$ | Crystal polarizability | DFT |
| $Stc$ | Stoichiometry | Formula |
| $Sym$ | Crystal symmetry | DFT |
| $a_m$ | Minimal lattice constant | DFT |
| $\Omega$ | Number of formula units | DFT |
| $\langle abc \rangle$ | Volume per formula unit | DFT |
| $\langle \alpha \beta \gamma \rangle$ | Standard deviation of lattice constants | DFT |
| $SO$ | Standard deviation of lattice angles | DFT |
| $N_{nn}$ | Number of crystal symmetry operations | DFT |
| $d_{nn}$ | Number of nearest neighbors | DFT |
| $d_{nn}$ | Nearest neighbors distance | DFT |
| $E_g$ | Band gap | DFT |
| $E_{vac}$ | Vacancy energy of the mobile ion | DFT |

**TABLE I.** Analyzed SSE descriptors and their abbreviations. The materials descriptors were classified into the categories (1) “mechanical and elastic” (M-E), (2) “diffusive and vibrational” (D-V) and (3) “structural and compositional” (S-C). The method of calculation of each descriptor, either zero-temperature (DFT) or finite-temperature (AIMD) simulations, is indicated in the third column. Some descriptors were directly deduced from the compounds formula, indicated as “Formula” in the table.
ACKNOWLEDGEMENTS

We acknowledge financial support from the MCIN/AEI/10.13039/501100011033 under Grant No. PID2020-119777GB-I00, the “Ramón y Cajal” fellowship RYC2018-024947-I, the Severo Ochoa Centres of Excellence Program (CEX2019-000917-S), the Generalitat de Catalunya under Grant No.2017SGR1506, and the CSIC under the “JAE Intro SOMdM 2021” grant program.

AUTHOR CONTRIBUTIONS

C.C. conceived the study and planned the research, which was discussed in depth with the rest of co-authors. C.C. and R.R. performed and analyzed the first-principles calculations. C.L. carried out the data analysis of the generated DFT-AIMD database as well as the training of the SSE machine learning models. A.E. created the website that gives access to the DFT-AIMD database. The manuscript was written by C.C. with substantial input from the rest of co-authors.

ADDITIONAL INFORMATION

Supplementary information is available in the online version of the paper.

COMPETING FINANCIAL INTERESTS

The authors declare no competing financial interests.