Investigations of the Optical Properties of GaNAs Alloys by First-Principle

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We present a Density Functional Theory (DFT) analysis of the optical properties of dilute-As GaN$_{1-x}$As$_x$ alloys with arsenic (As) content ranging from 0% up to 12.5%. The real and imaginary parts of the dielectric function are investigated, and the results are compared to experimental and theoretical values for GaN. The analysis extends to present the complex refractive index and the normal-incidence reflectivity. The refractive index difference between GaN and GaNAs alloys can be engineered to be up to ~0.35 in the visible regime by inserting relatively low amounts of As-content into the GaN system. Thus, the analysis elucidates on the birefringence of the dilute-As GaNAs alloys and comparison to other experimentally characterized III-nitride systems is drawn. Our findings indicate the potential of GaNAs alloys for III-nitride based waveguide and photonic circuit design applications.

III-nitride semiconductors have been extensively studied in the past decades and have been in the epicenter of semiconductor research in recent years1–4. The binary (Al, Ga, In)N alloys and their ternary and quaternary common-anion alloys possess remarkable optoelectronics, chemical, and tribological properties5,6 leading to a wide range of potential technological applications including that of solid state lighting, power electronics and MEMS devices7–9. Specifically, the invention of violet and blue GaN-based lighting emitting diode (LED) technologies has been recognized by the Nobel Prize in Physics in 201410.

The immense attention directed towards understanding and improving the conventional III-nitride alloys has led to significant breakthrough in the past decade. In comparison to GaN, InGaN and AlGaN alloys, the dilute-anion III-nitride-based material systems have received much less attention in the past decades11,12. In other words, the development in the dilute-anion III-nitride semiconductor class is still in the early stage and the understanding of these materials is still at its infancy. One of the prime examples in the dilute-anion III-nitride semiconductor class is the GaNAs with minute amount of arsenic (As) ~1–5%, which is called as “dilute-As GaNAs”; it is important to note that, its counterpart dilute-nitride InGaAsN alloy has been extensively investigated, which had resulted in state-of-the-art low threshold laser for telecommunication applications13–16.

The dilute-As GaNAs alloys have been experimentally realized via metal-organic chemical vapor deposition (MOCVD) with As-content reaching up to ~7%17,18. The GaNAs alloys with the whole As-content range have also been synthesized using molecular beam epitaxy (MBE) techniques19. While the experimental work demonstrated the feasibility to synthesize dilute-As GaNAs alloys, the physical properties of the alloy are still relatively unexplored. Recent work on the dilute-As GaNAs alloy shed light on its electronic properties, pointing out the large band gap coverage from ~2.3 eV to ~3.4 eV with As-content as little as 6%11,20–23. In addition, the interband Auger recombination process was suppressed within the dilute-As GaNAs alloy, indicating possible reduction of efficiency droop issue known in the InGaN based LEDs24. Thus, a novel design based on the InGaN/GaNAs interface quantum well (QW) has been proposed and it has been shown that the spontaneous emission rate for red emission can be enhanced by about 8.5 times compared to the conventional InGaN QW design25. These studies indicate a strong potential for implementing the dilute-As GaNAs semiconductor as the active region material in light-emitting devices.

The investigations of the electronic properties of dilute-As GaNAs alloy have provided a necessary foundation and possible ways for future device implementation. However, the other material properties of dilute-As GaNAs alloys, specifically the optical properties, are yet to be understood. Among the optical properties for example, the refractive indices of semiconductors have played a key role in photonics applications, including that of distributed

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Bragg reflector (DBR), waveguides and photonics crystals\textsuperscript{26–28}. Therefore, intensive understanding of the behavior of the optical properties of the dilute-As GaN\textsubscript{1-x}As\textsubscript{x} material is of high importance for future photonic device design.

In this work, we present an analysis of the optical properties of dilute-As GaN\textsubscript{1-x}As\textsubscript{x} alloys by employing First-Principle Density Functional Theory (DFT) calculations along with the use of a scissor operator. The arsenic (As) atoms are introduced as a replacement of nitrogen (N) atoms in the bulk GaN material system, in which the As-content is varied from 0% up to 12.5% As-content. The imaginary and real parts of the dielectric function of dilute-As GaN\textsubscript{1-x}As\textsubscript{x}, with comparison to theoretical and experimental values of GaN are specifically presented and discussed. Moreover, the complex refractive index, normal-incidence reflectivity, as well as birefringence, of the wurtzite material system are presented and comparison is drawn to other III-nitride material systems. Thus, we found that the changes in the refractive index of the GaN\textsubscript{1-x}As\textsubscript{x} alloys as compared to the GaN system can be significant, considering that the addition of arsenic is minute in the GaN material.

Computational Method
The DFT calculations in this study utilized the supercell approach to build the appropriate crystal structures of the dilute-As GaN\textsubscript{1-x}As\textsubscript{x} alloys to calculate the optical properties. In Fig. 1, a 128-atom supercell, comprised of 64 gallium (Ga) atoms, 63 nitrogen atoms and 1 arsenic atom is shown. By varying the supercell size and keeping one As-atom in the system at a time, the As-content in the alloys was varied, where the 128-atom supercell corresponds to 1.56% As-content. The projector augmented wave (PAW) method within the Vienna ab initio simulation package (VASP) was used to obtain the optical properties of the GaN\textsubscript{1-x}As\textsubscript{x} alloys\textsuperscript{29,30}. Moreover, the DFT calculations utilized the Generalized Gradient Approximation (GGA) and the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional for optimizing the structure\textsuperscript{31}. Thus, the optical properties were then calculated using the Local Density Approximation (LDA), where a scissor operator has been applied to shift the imaginary part of the dielectric function, similar to previous works\textsuperscript{32}. The cut-off energy was set to 400 eV and the external stress applied to the system was set to 0 GPa. Structure optimization was done using a Hellman–Feynman force of 0.02 eV/Å and an energy convergence criterion of $1 \times 10^{-5}$ eV/atom. Our calculations are performed based on the blocked Davidson algorithm while the tetrahedron smearing method was used as the integration scheme for the optical spectra. In addition, the number of bands and k-mesh sizes vary depending on the supercell size. As an example, different k-meshes were generated where the k-mesh for 128-atom supercell had a $3 \times 3 \times 3$ mesh, the 64-atom had a $7 \times 7 \times 7$ mesh and the 16-atom supercell had a $9 \times 9 \times 9$ mesh. Furthermore, we performed calculations including the spin-orbit coupling effect, to determine the effect of arsenic, but the results showed a negligible effect, similar to other wide band-gap III-nitride semiconductors and therefore has been excluded from the calculations in this work.

Complex Dielectric Function for Dilute-As GaN\textsubscript{1-x}As\textsubscript{x} Alloys. The complex dielectric function is defined as $\varepsilon(E) = \varepsilon_r(E) + i\varepsilon_i(E)$, where $\varepsilon_r(E)$ and $\varepsilon_i(E)$ are the real and imaginary parts of the complex dielectric function, respectively. The real and imaginary parts of the dielectric function are related by the Kramers-Kronig transformation in the following way\textsuperscript{33}:

$$\varepsilon_r(E) = 1 + \frac{2}{\pi} \int_0^\infty \frac{E'\varepsilon_i(E')}{E'^2 - E^2} dE'$$ (1)

$$\varepsilon_i(E) = -\frac{2E}{\pi} \int_0^\infty \frac{\varepsilon_i(E'')}{E'^2 - E^2} dE''$$ (2)

Therefore, once the imaginary part is obtained, it is relatively straightforward to obtain the real part of the dielectric function and vice versa.
Figure 2(a) and (b) show the imaginary parts of the complex dielectric function for dilute-As GaNAs alloys with As-content ranging from 0% up to 12.5%, for the xy- and z-directions, respectively. Moreover, the xy-direction is defined as the polarization vector perpendicular \((E_\perp c)\) to the surface of the GaNAs alloy, whereas the z-direction corresponds case parallel \((E_\parallel c)\) to the surface (c-axis) of the GaNAs alloy. Hence, the dilute-As GaNAs alloys exhibit anisotropy, similar to GaN and other ternary III-nitride material systems such as InGaN and AlGaN\(^{34,35}\).

From Fig. 2(a), increasing the arsenic content in the dilute-As GaNAs alloys leads to a shift in the onsets of the imaginary parts of the dielectric functions in the xy-direction towards lower photon energies. This phenomenon is analogous to the decrease in the energy band gap, which has been reported previously in the literature\(^{11,36}\). In addition, the positions of the E1, E2 and E3 absorption peaks of GaN in Fig. 2(a) can be observed at approximately 7.8 eV, 8.3 eV and 9.1 eV, respectively. The locations, as well as the magnitudes of the peaks compare well to those previously reported by Benedict and Shirley (LDA), whose peaks were around 7.5 eV, 7.9 eV, and 9 eV\(^{37}\). Similarly, Laskowski and co-workers (GGA) obtained peaks located around 7.5 eV, 7.9 eV and 9 eV for the E1, E2 and E3 peaks, respectively\(^{38}\). Moreover, increasing the arsenic content shifts the main absorption peaks of \(\varepsilon_{2,xy}(E)\) of the dilute-As GaNAs alloys towards the lower energy spectrum, with the 12.5% As-containing alloy showing the most prominent change in the overall shape of the \(\varepsilon_{2,xy}(E)\) spectra.

Figure 3(a) and (b) show the calculated real part of the dielectric function for dilute-As GaNAs alloys with As-content ranging from 0% up to 12.5%, for the xy- and z-directions, respectively. The real part of the complex dielectric function was obtained by using equation (1), as mentioned above. The values of \(\varepsilon_{1,xy}(E)\) for GaN are in good agreement with theoretical works from Persson and co-workers, as well as the experimental and theoretical works for GaAs and InAs. The agreement on the imaginary part of the dielectric function for GaN with previous works provides a strong confidence on the method we used in this work to obtain the optical properties of the dilute-As GaNAs alloys. Moreover, increasing arsenic content in the dilute-As GaNAs alloys has more prominent changes in the magnitude of \(\varepsilon_{2,xy}(E)\) spectra (Fig. 2(b)) compared to the perpendicular polarized case (Fig. 2(a)), up to the 6 eV region. Thus, the shift in the main absorption peaks (E1-E3) of \(\varepsilon_{2,xy}(E)\) spectra for GaNAs alloys containing 1.56% arsenic up to 12.5% are incremental with respect to DFT-calculated GaN.

Figure 3(a) and (b) show the calculated real part of the dielectric function for dilute-As GaNAs alloys, with As-content ranging from 0% up to 12.5%, for the xy- and z-directions, respectively. Note that the real part of the complex dielectric function was obtained by using equation (1), as mentioned above. The values of \(\varepsilon_{1,xy}(E)\) for GaN are in good agreement with theoretical works from Persson and co-workers, as well as the experimental and theoretical works for GaAs and InAs. The agreement on the imaginary part of the dielectric function for GaN with previous works provides a strong confidence on the method we used in this work to obtain the optical properties of the dilute-As GaNAs alloys. Moreover, increasing arsenic content in the dilute-As GaNAs alloys has more prominent changes in the magnitude of \(\varepsilon_{2,xy}(E)\) spectra (Fig. 2(b)) compared to the perpendicular polarized case (Fig. 2(a)), up to the 6 eV region. Thus, the shift in the main absorption peaks (E1-E3) of \(\varepsilon_{2,xy}(E)\) spectra for GaNAs alloys containing 1.56% arsenic up to 12.5% are incremental with respect to DFT-calculated GaN.

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Figure 3 shows that the high-frequency dielectric function for GaN was deduced to be \(\varepsilon_{\infty}(E)\approx4.8\), which compares well to previously reported theoretical values from Lambrecht and co-workers (~4.79), but is slightly lower compared to Takeuchi and co-workers (~5.27) and de Carvalho and co-workers (~5.11)\(^{31-43}\). Thus, Fig. 3(a) indicates that increasing the As-content in the dilute-As GaNAs alloys leads to higher values of the high-frequency dielectric constant \(\varepsilon_{xy}(\infty)\), with the 12.5% As-content alloy having a value of \(\varepsilon_{xy}(\infty)\approx5.85\). Additionally, the static dielectric constant \(\varepsilon_{xy}\) for DFT-calculated GaN is \(\approx8.7\), which is in good agreement with experimental results found by Strite and Morkoç (~8.9)\(^{44}\). On the other hand, increasing the arsenic content leads to slightly higher values in the static dielectric constant, with the 12.5% As-containing alloy having \(\varepsilon_{xy}\approx9\). Although these results have not yet been experimentally confirmed, they can be useful in simulation modelling for future photonic device applications.

On the other hand, Fig. 3(b) indicates that the \(E \parallel c\) polarized real part of the dielectric function of our DFT-calculated results is in good agreement with those of Benedict and co-workers\(^{45,46}\). The DFT-calculated GaN high-frequency dielectric constant for the \(z\)-direction resulted in a value of \(\varepsilon_{z}(\infty)\approx4.95\) - which is slightly higher than that reported by Christensen \textit{et al}. (~4.62) and slightly lower compared to de Carvalho \textit{et al}. (~5.3)\(^{31,43}\). The discrepancy can be ascribed to the difference in the method for calculating the \(\varepsilon_{z}(E)\) spectra, as Christensen \textit{et al}. did not include any LDA correction terms. Thus, compared to Fig. 3(a), Fig. 3(b) shows that the magnitudes of the dielectric functions in the \(z\)-direction are slightly higher for all the presented cases as compared to the \(xy\)-direction. This trend, which is directly related to the anisotropic behavior of wurtzite material systems, has been observed experimentally in various works including those of Shokhovets, Pezzagna, Rigler with co-workers, and theoretically by Christensen \textit{et al}\(^{45-48}\). Thus, Fig. 3 shows that the high-frequency dielectric function for all \(\varepsilon_{z}(\infty)\) values is higher compared to \(\varepsilon_{xy}(\infty)\) for all dilute-As GaNAs alloy compositions studied, and that the increase is a linear relationship of arsenic content. This indicates a possibility allowing for precise control of the optical properties of the GaN-based alloy by simply tuning the arsenic content.

**Complex Refractive Index, Reflectivity and Birefringence of Dilute-As GaNAs Alloys.** By using the information obtained from the complex dielectric function, optical constants like the complex refractive index \(n'(E) = n(E) + ik(E)\) can be obtained with the following relations\(^{33}\):

\[
n(E) = \sqrt{\frac{\varepsilon_1(E)^2 + \varepsilon_2(E)^2 + \varepsilon_3(E)^2}{2}} \tag{3}
\]

\[
k(E) = \sqrt{\frac{\varepsilon_1(E)^2 + \varepsilon_2(E)^2 - \varepsilon_3(E)^2}{2}} \tag{4}
\]

Figure 4(a) and (b) show the refractive index for the \(xy\)- \((n_{xy}(E))\) and \(z\)-directions \((n_{z}(E))\) for DFT-calculated dilute-As GaNAs alloys, with As-content varying from 0% up to 12.5%, along with experimentally fitted data for GaN, respectively. Figure 4(a) indicates the refractive index in the \(xy\)-direction increasing linearly as the arsenic content in the dilute-As GaNAs alloys is increased for photon energies up to ~3.5 eV. Moreover, the E0 peak (band edge) of DFT-calculated GaN is approximately located around ~3.4 eV, which agrees well with data reported by Takeuchi and Djurišić with co-workers\(^{41,49}\). The respective E0 peaks for dilute-As GaNAs alloys shift to lower photon energies as more arsenic is introduced into the system, with the 12.5% As-containing alloy having the E0 peak around ~2.2 eV. Furthermore, Fig. 4(a) shows that the DFT-calculated refractive index of GaN \((n_{GaN})\) is about ~5% on average lower (below the band gap) compared to experimentally measured values of Shokhovets and co-workers\(^{40}\). It should be noted that our calculations were performed at 0 K and that effects like strain and surface roughness were not taken into account, albeit previously identified to affect the magnitude of the dielectric function\(^{42,45}\). In addition, Rigler and co-workers noted that Ga-polar and N-polar thin films have different

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**Figure 4.** (a) Refractive index in the \(xy\)-direction \((n_{xy}(E))\) and (b) \(z\)-direction \((n_{z}(E))\) spectra for dilute-As GaNAs alloys (solid lines), with As-content ranging from 0% up to 12.5% arsenic, and experimentally fitted data (dashed lines), respectively.
refractive indices, which they ascribed it to an effect due to different impurity contents. The effects of temperature on the refractive index of dilute-As GaNAs alloys will be essential for future photonic device design, but this study is not within the scope of this work and will require further investigations.

Figure 4(b) indicates that the experimental values of $n(E)$ from Shokhovets and co-workers for GaN are about 5% higher on average compared to the values of DFT-calculated GaN, in the energy range from ~0.5 eV up to ~3.3 eV. Similar to the xy-direction, the E0 peaks of the refractive indices in the z-direction of dilute-As GaNAs alloys shifted to the left, as more arsenic is introduced into the GaN-based system, with the 12.5% As-containing alloy showing a peak around ~2.3 eV. Thus, by inserting a minute amount of arsenic into the GaN material system, it is possible to alter the material system in such a way to achieve controlled large refractive index differences compared to GaN. On the other hand, in a material system like InGaN, it has been reported that the refractive index is strongly dependent on numerous factors including phase-separation, the piezoelectric-field induced Stark-effect and variations due to quantum-confinement, making it difficult to predict the refractive index based solely on the indium content. Furthermore, AlGaN has a similar behavior to dilute-As GaNAs alloys, where the refractive index is a linearly decreasing function for increasing Al-content. It should be noted that Fig. 4(a) and (b) do not indicate the alloy compositions that may be absorbing light, for example, the GaN$_{0.875}$As$_{0.125}$ alloy will be absorbing the 450 nm light, since the bandgap is around ~2.23 eV.

Figure 5(a) and (b) present the imaginary part of the complex refractive index data in the xy- and z-directions for dilute-As GaNAs alloys with As-content varying from 0% up to 12.5%, respectively. For the $k_x(E)$ spectra in Fig. 5(a), the absorption onset is noticeable around ~3.5 eV, which is consistent with the reported values of the energy band gap for GaN. Moreover, the overall shape and magnitude of our DFT-calculated and the GaN data reported by Lambrecht et al. are in good agreement for energies up to 10 eV. Furthermore, increasing the As-content in dilute-As GaNAs alloys shifts the $k_x(E)$ spectra towards the lower photon energy region, consistent with the observed energy gap reduction from previous works. Thus, in Fig. 5(a), the overall magnitude of $k_x(E)$ increases as more arsenic is introduced into the dilute-As GaNAs system, which is not the case for the data presented in Fig. 5(b). Namely, the 12.5% As-containing GaNAs alloy has a noticeably lower $k_x(E)$ value in the range from 7 eV to 10 eV compared to all other systems considered. In addition, the overall shape, of $k_x(E)$ and $k_z(E)$ of DFT-calculated GaN coincides well with previous theoretical works of Lambrecht and co-workers, as well as Djurišić and Takeuchi with co-workers, although the magnitude is slightly lower in our case as compared to the latter two.

By using the results on the complex refractive index, it is possible to extract the normal-incidence reflectivity for the dilute-As GaNAs alloys by employing the following relation:

$$R(E) = \frac{\left|n(E) - 1\right|^2 + k(E)^2}{\left|n(E) + 1\right|^2 + k(E)^2}$$  \hspace{1cm} (5)

Figure 6(a) and (b) present the normal-incidence reflectivity for dilute-As GaNAs alloys with As-contents ranging from 0% up to 12.5%, for the xy- and z-directions, respectively. From Fig. 6(a), the reflectivity for DFT-calculated GaN in the xy-direction ($R_x(E)$) falls in the range of ~0.13–0.18 for energy ranging from 0 eV to 3 eV, which is in good agreement with the values reported by Kawashima and co-workers. Similarly, Fig. 6(b) also indicates that the normal-incidence reflectivity for DFT-calculated GaN in the z-direction ($R_z(E)$) lies in the range of 0.14–0.18. On the other hand, the 12.5% As-containing GaNAs alloy has the highest $R_x(E)$ and is varying between ~0.17–0.2, within the visible regime. Thus, when the arsenic content in the dilute-As GaNAs alloys is increased, a linear enhancement in the normal-incidence reflectivity for the xy-, as well as z-directions is observed, as shown in Fig. 6(a) and (b), for energies up to ~3 eV.

The birefringence of the dilute-As GaNAs alloy can be extracted out by comparing its real part of the complex refractive indices from both directions (xy and z). The birefringence is defined as $\Delta n = n_x - n_y$, where $n_x$ is the refractive index in the z-direction and $n_y$ is the refractive index in the xy-direction. Figure 7 presents the birefringence ($\Delta n$) as a function photon energy of dilute-As GaN$_{1-x}$As$_x$ alloys, where the arsenic content varies from 0% up to 12.5%. The birefringence of DFT-calculated GaN aligns well with previously reported experimental values where a value of $\Delta n \sim 0.0425$ by Pezzagna and co-workers was obtained, whereas Sanford and co-workers...
Hui and co-workers reported $\Delta n \sim 0.04$ and Shokhovets and co-workers have reported a varying birefringence, that increases with increasing photon energy and ranges between 0.037 and 0.056. In this work, we obtained a birefringence for GaN that increases slowly from 0.044 to 0.051 for energies approaching the band edge. In addition, Rigler and co-workers have reported that birefringence may vary depending whether the GaN thin films are Ga-polar or N-polar, with the Ga-polar exhibiting slightly higher refractive indices – something that has not been taken into account in this work. On the other hand, it is important to notice that the birefringence of the dilute-As GaNAs alloys can be as low as ~0.02 when the As-content is relatively minute (1~2%). However, while approaching the bandgap value (~3.3 eV) for GaNAs with 1.56% As-content, its birefringence increases and exceeded that of GaN. This feature may prove interesting, as it may indicate a potential pathway to reduce the birefringence in the nitride-based system by simply inserting low amounts of arsenic into the GaN system. Additional work is necessary to address this reduction and determine the influence of temperature on the birefringence, as well as the refractive index. In general, our results indicate an increasing trend in $\Delta n$ for higher As-content in the dilute-As GaNAs alloy.

**Conclusion**

In summary, First-Principle DFT calculations were carried out to evaluate the optical properties of dilute-As GaNAs semiconductor alloys with As-content ranging from 0% to 12.5%. By using the knowledge on the complex dielectric function, optical constants including the complex refractive index and normal-incidence reflectivity were evaluated and analyzed. Our findings showed that dilute-As GaNAs alloys are optically anisotropic, and the analysis on the birefringence of the alloys was compared to other conventional III-nitride material systems. In addition, our analysis indicated large index difference (up to ~0.35) between dilute-As GaNAs alloy and GaN alloy across a large range of photon energies. Specifically, the refractive index of dilute-As GaNAs had a linear relationship with respect to the arsenic content in the alloys, indicating potential increment at even higher arsenic contents. The large refractive index difference tuning capability between GaN and dilute-As GaNAs alloys could prove beneficial for the design of waveguides and distributed Bragg reflectors, requiring the use of large refractive index differences. Our work points out the promising optical properties of dilute-As GaNAs semiconductors and pose them as strong candidates for materials for photonic device applications.

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Author Contributions
D.B., C.K.T., and N. T. contributed to the discussions, concept development, theoretical analysis, analysis of the results, and writing of the manuscript. N.T. supervised the studies performed in the manuscript.

Additional Information
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