The Effect of Bars on the Ionized ISM: Optical Emission Lines from Milky Way Analogs

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Abstract

Gas interior to the bar of the Milky Way has recently been shown to be the closest example of a low-ionization (nuclear) emission region (LI(N)ER) in the universe. To better understand the nature of this gas, a sample of face-on galaxies with integral field spectroscopy is used to study the ionized gas conditions of 240 barred and 250 nonbarred galaxies, focusing on those that are most similar to the Milky Way. Strong optical line emission of [N II] λ6584, Hα, [O III] λ5007, and Hβ are used to diagnose the dominant ionization mechanisms of gas across galaxies and the Galaxy via Baldwin–Phillips–Terlevich diagrams. Barred galaxies show a strong suppression of star formation and an increase in composite and LI(N)ER-like spectra in their inner regions when compared with similar nonbarred counterparts. This effect is lessened in galaxies of very low (log(M*/M⊙) ≤ 10.4) or very high (log(M*/M⊙) ≥ 11.1) total stellar mass. Bar masks from Galaxy Zoo:3D show the bar’s nonaxisymmetric effect on the ionized gas and help predict the face-on distribution of ionized gas conditions near the bar of the Milky Way.

Unified Astronomy Thesaurus concepts: Barred spiral galaxies (136); Milky Way Galaxy (1054); LINER galaxies (925); Interstellar line emission (844); Interstellar atomic gas (833); Interstellar radiation field (852); Interstellar medium (847)

Supporting material: animations

1. Introduction

Optical line emission was first detected in an extragalactic system in 1909 (Fathi 1909), although its extragalactic nature was not posited until 1929, when Cepheids were observed in M31 (Hubble 1929b). The correlation between distance and recession velocity (Hubble 1929a), the Hubble–Lemaître law, sparked large spectroscopic campaigns of galaxy observations (see Humason et al. 1956). In particular, a coordinated effort at the Mount Wilson and Lick Observatories started in 1935 and showed many galaxies with [O II] λ3727 Å nebular emission in their central regions (Humason 1936; Mayall 1939; Baade & Mayall 1951). Around the same time, Babcock (1939) tentatively detected [O II] λ3727 Å emission from the inner ~1 kpc of M31 (later confirmed by Münch 1960), where no star formation is present. Samples of galaxies showing strong optical emission lines from ionized gas continued to rise, and soon Seyfert (1943) discovered bright, highly ionized gas in the centers of some galaxies. Burbidge & Burbidge (1962, 1965) confirmed and extended these initial ionized gas detections using Hα and [N II] λ6584 emission across 85 galaxies and found that the line ratio of [N II]/Hα increased from 1/3 in outer regions and spiral arms to at or above unity in many galactic nuclei. Ratios this large were not seen within the Milky Way either at that time (Courtoës & Cruvellier 1960) or for the next half century (Haffner et al. 2009). Burbidge & Burbidge (1965) noted that such large line ratios could only be explained with an enhanced nitrogen abundance or a large electron temperature of 10,000–20,000 K (Burbidge et al. 1963). Over time, additional optical emission lines were studied and larger samples of galaxies were categorized based on their line ratios and underlying sources of ionization (Heckman 1980; Baldwin et al. 1981).

Burbidge–Phillips–Terlevich (BPT; Baldwin et al. 1981) diagrams use a combination of the [N II]/Hα and [O III]/Hβ line ratios to classify extragalactic systems. This method can identify the dominant source of excitation in galaxies using only line emission, allowing for source classification in large surveys. Veilleux & Osterbrock (1987) improved on this method, and many theoretical and empirical studies of large galaxy samples yielded four broad categories based on the source of ionization: star formation, active galactic nuclei (AGN), composites, and low-ionization nuclear emission.
regions (LINERs; Kewley et al. 2001; Kauffmann et al. 2003; Schawinski et al. 2007). While star formation and AGN have known sources of ionization, LINERs are a descriptive name for an unknown source first coined by Heckman (1980). In this work, we place the term “nuclear” in parentheses (Li(N)ER), since in many cases, this class of emission is seen outside of galactic nuclei (e.g., Sharp & Bland-Hawthorn 2010; Annibali et al. 2010; Belfiore et al. 2016). Some Li(N)ERs may arise from low-luminosity AGN, while others are correlated with older stellar populations (Kewley et al. 2006; Belfiore et al. 2016). The characteristic signatures of Li(N)ERs are lower values for [O III]/Hβ than in AGN but higher ratios for lower-ionization lines, such as [N II]/Hα, [S II]/Hα, and [O I]/Hα (Ho et al. 1993, 1997; Filippenko 2003). The large [N II]/Hα line ratios require a significant source of ionizing radiation in regions of galaxies where little to no star formation is occurring, as Babcock (1939) observed for M31.

Until recently, Li(N)ER gas was only observable in extragalactic systems, with the closest such example being M31. Optical observations of the ionized gas counterpart to the neutral Milky Way structure, the tilted disk (Burton & Liszt 1978; Liszt & Burton 1978, 1980), have allowed for the inner Milky Way to be classified using the same diagnostic method of BPT diagrams (Krishnarao et al. 2020). The Milky Way has been shown to have a bar (see Bland-Hawthorn & Gerhard 2016), most recently as a stellar density enhancement using Gaia, Pan-STARRS1, 2MASS, and AllWISE data (Anders et al. 2019). Gas interior to the bar around 0.5–1.5 kpc from Galactic center is largely ionized with optical emission line ratios similar to Li(N)ERs (Krishnarao et al. 2020). This gas is now the closest Li(N)ER to us in the universe at a distance of only ∼7 kpc. While the rest of the Milky Way does not yet have widespread measurements of all four emission lines used for the BPT diagram, it is rare locally and in spiral arms to have [N II]/Hα line ratios as large as those seen in Li(N)ERs (see Haffner et al. 2009, for a review). This suggests that Li(N)ER gas in the Milky Way is constrained to the vicinity of the bar and raises the question of how important a bar is in setting the ionized gas conditions and radiation field.

Simulations of gas in barred galaxies predict strong shear and noncircular flows across many phases of gas (Renaud et al. 2013; Seo et al. 2019; Sormani & Barnes 2019). They also predict concentrated locations of star formation in the form of rings, along with bursty star formation in a central molecular zone (CMZ)–like core. Such ringlike structures have been seen in optical images (e.g., Buta 1986) and are starting to be studied in nearby barred galaxy observations with the MUSE integral field unit (IFU; see Gadotti et al. 2019, and the TIMER survey). Galaxies with redder colors and stronger bulges have been shown to have larger bar fractions (Masters et al. 2011). Stars in the bar also show flatter age and metallicity gradients than stars in the disk at similar radii (Fraser-McKelvie et al. 2019), suggesting a large degree of radial mixing in the bar. Barred galaxies also have a lower atomic gas content, potentially inhibiting star formation in their host galaxy (Masters et al. 2012; Newholm et al. 2019). In addition, Hoyle et al. (2011) found that longer bars are redder and are present in redder galaxies, establishing a strong link between star formation suppression and bars. George et al. (2019) showed that in M95, star formation quenching in the vicinity of the bar is consistent with strong gas inflows that leave the bar devoid of neutral and molecular hydrogen that can form stars.

However, even with this suppression of star formation, atomic ionized gas is still present and often displays enhanced optical line ratios similar to Li(N)ERs.

Since the bar provides a nonaxisymmetric potential, gas is expected to transition from circular rotation to noncircular rotation near the bar radius. The effect of the resulting shocks associated with this transition and gas flows in a bar potential should be seen as changes in the BPT classifications of gas within galaxies. Kinematic models are required to analyze the distribution of gas in the inner Milky Way, where radial velocities must be mapped to distances. While some models of this nature are available, such as with the tilted disk in the inner Milky Way (Krishnarao et al. 2020), a complete picture cannot be seen from within the Galaxy. In particular, the behavior of the interstellar medium (ISM) interior to the bar radius but outside of the bar itself is difficult to ascertain with direct observations in our Galaxy. Face-on extragalactic systems provide a better way to obtain a more complete view of the ISM in the vicinity of a bar, both interior to and outside of its nonaxisymmetric distribution.

Studies focusing on extragalactic Milky Way analogs help place the Milky Way on extragalactic scaling relations (e.g., Robotham et al. 2012; Geha et al. 2017; Boardman et al. 2020; Fraser-McKelvie et al. 2019). These analog samples have been used to infer statistical properties of Milky Way–like galaxies and explore the extent to which such systems are outliers from global galaxy relations (Licquia et al. 2015; Licquia & Newman 2015; Boardman et al. 2020). They are selected based on stellar mass, star formation rate (SFR), bulge prominence, or other global properties but not on the presence of finer morphological structures, such as bars. Since a bar can have large-scale impacts on the ISM and star formation, it could be a vital characteristic of a galaxy to consider when constructing a sample of Milky Way analogs. In this work, we focus on better understanding Li(N)ERs and determining how crucial the presence of a bar is for the ionized gas properties of galaxies using SDSS MaNGA.

This paper is organized as follows. In Section 2, we describe the SDSS MaNGA observations and additional data products used, and we explain our sample selection process in Section 3. Section 4 presents our comparisons of barred and nonbarred galaxies as a function of various galaxy properties, and Section 5 focuses on the barred sample to reveal the nonaxisymmetric effects on the ISM of galaxies. We discuss the broader implications of our findings in Section 6 in the context of the Milky Way and close with our conclusions in Section 7.

2. Data

This paper uses data and products associated with the Sloan Digital Sky Survey (SDSS; York et al. 2000). The details of these data products are described below.

2.1. MaNGA IFU Data

The Mapping Nearby Galaxies at Apache Point Observatory (MaNGA) survey is a part of SDSS-IV (Blanton et al. 2017) that uses integral field spectroscopy to observe 3D spectra of nearly 10,000 galaxies (Bundy et al. 2015). Spectra are obtained with the BOSS spectrograph on the 2.5 m telescope at Apache Point Observatory (Gunn et al. 2006) at a spectral resolution of $R \sim 2000$ for 3600 Å < λ < 10300 Å and have
exposure times such that a target signal-to-noise ($S/N$) level is achieved (Bundy et al. 2015). Fibers subtend 2′′ on the sky (Smee et al. 2013) and are bundled into IFUs with sizes 12″–32″ in diameter with 19–127 fibers (Drory et al. 2015). Observed data undergo sky subtraction and flux calibration using simultaneous observations of the sky and standard stars (Yan et al. 2016). The median point-spread function for MaNGA data cubes is 2.5″ and roughly corresponds to kiloparsec physical scales at the targeted redshift range (0.01 < $z$ < 0.15). Observations are dithered and mapped onto spectroscopic pixels (or spaxels) that are 0.5″ across.

MaNGA targets local galaxies, and the sample is selected to have a flat distribution of $i$-band absolute magnitude and uniform radial coverage. This sample is composed of three main components, a primary sample where 80% of galaxies are covered out to 1.5 $R_e$, a secondary sample where 80% of galaxies are covered out to 2.5 $R_e$, and a color-enhanced supplement to improve coverage of poorly sampled regions of the NUV – $i$ versus $M_i$ color–magnitude plane (Wake et al. 2017). All MaNGA data in this work are reduced using v2.5.3 of the MaNGA Data Reduction Pipeline (DRP; Law et al. 2016) and employ the Data Analysis Pipeline (DAP; Belfiore et al. 2019; Westfall et al. 2019) from the internal eighth MaNGA product launch (MLP-8) that contains observations of 6507 galaxies. The DAP provides emission line data products that this work employs ([N II] $\lambda$6584, H$\alpha$, [O III] $\lambda$5007, and H$\beta$) and can be accessed using the Python package sdss-marvin (Cherinka et al. 2019). The processed data files used here follow a hybrid binning scheme where spectra are first Voronoi-binned (Cappellari & Copin 2003) to reach a $g$-band $S/N$ ~10 for determining stellar kinematics, while emission lines are fit to each spaxel within the bin. Throughout this work, we assume a standard cosmology of WMAP9 (Hinshaw et al. 2013).

2.2. Galaxy Zoo 2

Morphological classifications for all SDSS MaNGA parent sample galaxies are from the citizen science Galaxy Zoo 2 project (Willett et al. 2013; Hart et al. 2016). Volunteer citizen scientists used SDSS galaxy images to identify galaxies as early types, late types, or mergers, along with measuring finer features such as bars, bulges, and shapes of edge-on disks. Classifications are made following a decision tree model, where users answer questions about a galaxy and receive follow-up questions based on their response. Willett et al. (2013) provided unbiased vote fractions, improved upon in Hart et al. (2016), for each of these questions, along with thresholds to determine well-sampled galaxies for each classification. In this work, the thresholds recommended in Willett et al. (2013, their Table 3) are used to find a barred and nonbarred galaxy sample.

2.3. Total Stellar Masses and SFRs

Total stellar masses for MaNGA galaxies are from the NASA-Sloan Atlas (NSA; Blanton et al. 2011) and derived from an elliptical Petrosian photometric analysis. Estimates of the total SFR are from the DAP (Belfiore et al. 2019; Westfall et al. 2019) and calculated using the Gaussian fitted H$\alpha$ flux within the IFU field of view following the work of Kennicutt & Evans (2012), Murphy et al. (2011), and Hao et al. (2011) and assuming a Kroupa initial mass function (IMF; Kroupa 2001). These SFR measurements do not remove non-star-forming spaxels (e.g., AGN or LI(N)ER BPT classifications) or correct for extinction. They provide an estimate of the star formation within the IFU footprint of a galaxy and are only used as a brief check.

2.4. Stellar Mass Surface Densities

Pace et al. (2019a) provided robust estimates of resolved stellar mass for all SDSS MaNGA galaxies from MPL-8 using a principal component analysis (PCA) method for a more computationally tractable stellar continuum fit. A set of six eigenspectra is defined by performing PCA on a set of 40,000 synthetic “training” spectra, themselves obtained from applying procedurally generated star formation histories to a theoretical library of stellar atmospheres (see Conroy & van Dokkum 2012) using the Kroupa (2001) IMF. The PCA basis set of six eigenspectra spans a lower-dimensional space, inside which the spectral fitting occurs, but the space expresses >99% of the variance in the training library. Each spectrum (whether from the training set or observed) can be expressed as a linear combination of the six eigenspectra,

$$S_i = Q \left( S_{\text{med}} + \sum_i A_{PC,i} \cdot S_{PC,i} \right)$$

where $Q$ is a normalization factor, $S_{\text{med}}$ is the median basis spectrum, $A_{PC,i}$ are the amplitudes of the principal component spectra, and $S_{PC,i}$ are the principal component eigenspectra. The SDSS $i$-band stellar mass-to-light ratio for a given spectrum $S_i$ is obtained by evaluating the likelihood of each training spectrum $S_{\lambda_{\text{lam}}}$ against the observation $S_i$ in the lower-dimensional principal component space. The corresponding values of the stellar mass-to-light ratio for each $S_{\lambda_{\text{lam}}}$ are then treated like samples from a prior and weighted according to their likelihood: the weighted sample median is taken as the best estimate of the stellar mass-to-light ratio. This method is reliable at a wide range of $S/N$s ($2 < S/N < 30$), at low-to-moderate dust attenuation ($\tau \lesssim 4$), and over the full range of realistic stellar metallicities.

The stellar mass-to-light ratios are converted to stellar mass surface density, first by computing and multiplying by the reconstructed $i$-band magnitude of the starlight from the principal component fit, then dividing by the projected area of a MaNGA spaxel at the galaxy’s redshift, and finally deprojecting using the minor-to-major axis ratio from the NSA (Blanton et al. 2011).

2.5. Galaxy Zoo:3D

Ongoing citizen science efforts associated with the Galaxy Zoo project provide masks to isolate emission from structural components of galaxies, including bars. These Galaxy Zoo:3D (Masters et al. 2020, in preparation) masks allow for emission originating within the bar to be compared with emission outside of the bar. Details of this ongoing project can be found on their website (https://www.zooniverse.org/projects/kgmasters/galaxy-zoo-3d), and preliminary data products have been used to study density wave theory (Peterken et al. 2019) and stellar population gradients (Fraser-McKelvie et al. 2019). While qualitative in nature, these masks effectively separate MaNGA spaxels dominated by bar light from the rest of the galaxy. Spaxels flagged to be in the bar by at least 20% of participants are considered to be in the
bar, as done in Fraser-McKelvie et al. (2019). Out of the 240 barred galaxies in our full MPL-8 sample (described in Section 3), 213 have reliable bar masks. The maximum radius of spaxels in a bar mask and the minimum radius of spaxels outside a bar mask provide measures of the length and width of the bar, respectively. These bar length measurements are validated with independent measures of the bar length using a Fourier decomposition method employed in Kraljic et al. (2012; see Appendix A for details).

3. Galaxy Sample

The ISM is known to be strongly affected by bars in galaxies, as evidenced by many N-body simulations (e.g. Athanassoula 1992; Debattista & Sellwood 1998) and observational studies (e.g., Hoyle et al. 2011; Masters et al. 2012; George et al. 2019). Because the Milky Way is a prominently barred galaxy (Binney et al. 1991; Weiland et al. 1994; Benjamin et al. 2005; Bland-Hawthorn & Gerhard 2016), it is crucial to understand the impact of bars on the ISM, but doing so from within the Galaxy is limited by our view from within. Instead, face-on galaxies allow us to examine varying ionized gas properties of the ISM as a result of the presence of bars. Four samples of galaxies are considered: barred Milky Way analogs, nonbarred analogs, a full barred galaxy sample, and a full nonbarred galaxy sample. The analog samples are constructed of galaxies in the MaNGA parent sample based on morphology, color, and stellar mass, while our full barred and nonbarred samples follow the same criteria but do not consider a stellar mass range. A mass selection is applied to nonbarred galaxies to match the stellar masses of barred galaxies.

The complete MaNGA survey will observe approximately 10,000 out of 29,811 galaxies from the parent sample. Our sample is initially constructed from all galaxies in this parent sample. This allows for our results to be easily replicated with additional galaxies as they are observed up until the completion of the MaNGA survey. A decision tree of the sample selection process is shown in Figure 1 and described in the following subsections.

3.1. Morphology

First, the morphology of the SDSS MaNGA parent sample of galaxies is restricted using Galaxy Zoo 2 (Willett et al. 2013) to those with a face-on disk. This process starts with selecting galaxies with at least 43% of votes identifying a disk ($P_{\text{features/disk}} > 0.43$), followed by at least 71.5% of votes identifying the galaxy as not edge-on ($P_{\text{not edge-on}} > 0.715$) with at least 20 such votes ($N_{\text{not edge-on}} \geq 20$). These face-on disk galaxies are then sorted into galaxies with and without a bar using at least 80% of votes identifying a bar ($P_{\text{bar}} \geq 0.8$; barred galaxies) and less than 20% of votes identifying a bar ($P_{\text{bar}} \leq 0.2$; nonbarred galaxies). This results in a total of 1380 barred galaxies and 3294 nonbarred galaxies in the SDSS MaNGA parent sample. These parameters are chosen based on the recommendations of Willett et al. (2013), and the $P_{\text{bar}} \geq 0.8$ parameter is a very tight constraint on the initial set of galaxies. This tight constraint reduces our completeness of barred galaxies to about 29.4% when compared with the visual classifications of Nair & Abraham (2010) but ensures a high accuracy in identifying true bar structures.

3.2. Star Formation and Mass Selection

To consider the ISM within our sample, we select galaxies that show signs of star formation using an NUV \(- r \leq 5$ color
cut, as done in Belfiore et al. (2017). We use this minimal cut rather than selecting galaxies based on SFR because of the difficulty in estimating an accurate SFR for all MaNGA galaxies and, similarly, the difficulty in estimating a global SFR of the Milky Way. The DAP estimates of the SFR do not account for extinction and take a simple approach to providing consistent measurements for all galaxies (Belfiore et al. 2019; Westfall et al. 2019). Within the Milky Way, systematic differences between different methods are still present (see, e.g., McKee & Williams 1997; Chomiuk & Povich 2011). This loose selection restricts our barred and nonbarred galaxies to 1262 and 3172, respectively. With these criteria, the barred and nonbarred galaxy samples show a noticeable difference in their total elliptical Petrosian stellar mass distributions from the NSA (Blanton et al. 2011), with nonbarred galaxies showing a greater presence at lower masses. To ensure our sample selection process has not inadvertently selected on other parameters, we use the $k$-sample Anderson–Darling statistical test to compare many parameter distributions from our samples. The $k$-sample Anderson–Darling statistical test (Scholz & Stephens 1987) compares the underlying populations for each sample with a null hypothesis that all samples come from the same population. We prefer the Anderson–Darling test to the more common Kolmogorov–Smirnov test for its more sensitive statistical properties (see Babu & Feigelson 2006). Ideally, the barred and nonbarred samples would fail to reject this null hypothesis for all galaxy parameters, indicating that our barred and nonbarred samples are probing the same population of galaxies. This ensures our study is not affected by other galaxy parameters that can impact the ISM, isolating the effects of the bar.

The stellar mass distributions for the parent barred and nonbarred samples allow for this null hypothesis to be rejected at the 0.01 level, suggesting that the barred and nonbarred samples come from distinct underlying populations. To avoid this stellar mass bias, a mass selection is applied to the nonbarred galaxies, where for each barred galaxy, the nonbarred galaxy closest in total stellar mass is chosen. Figure 2 shows the results of this process, and our mass-selected barred and nonbarred samples show a statistically similar distribution, with the $k$-sample Anderson–Darling statistical test $p$-value $>0.25$. Similar tests comparing the SFR (from the DAP; Westfall et al. 2019), far-UV (FUV), near-UV (NUV), and $i$-band photometry (from the NSA; Blanton et al. 2011) of these samples also fail to reject the null hypothesis at the $p = 0.01$ level, indicating that the samples are drawn from similar populations. In this work, we use observations available as of MPL-8, which results in our full MPL-8 barred and nonbarred galaxy samples with sizes of $N = 240$ and 250, respectively. The bottom panel of Figure 2 shows the total stellar mass distributions and statistical test results for the full MPL-8 samples. While small differences are present in these distributions, they are mostly similar ($p$-value $>0.01$) and should provide a fair comparison of barred versus nonbarred galaxies.

The presence of a different underlying distribution of stellar mass across our barred and nonbarred samples may introduce another source of bias. Differences in the ISM of our galaxy samples may be influenced by differences in the bulge properties of our two samples. To test this, we consider the stellar mass surface density across each galaxy for our barred and nonbarred samples. In Figure 3, a violin-style Gaussian kernel density histogram of the stellar mass surface density per spaxel from the PCA analysis of Pace et al. (2019a) is shown along different effective radius bins for barred and nonbarred galaxies. This style of plot is used throughout this paper to compare the distributions of properties across categories. The shapes of the violin plots are from the probability density distribution, where widths are normalized by the number of spaxels in each category. The relative widths of the violin shapes across different categories at a given bin...
value. While small differences can be seen in the inner density at different radial bins of barred galaxies, corresponding to the fraction of spaxels within a category at that bin value. In general, both barred and nonbarred galaxies should primarily result from the impact of the bar.

4. Barred versus Nonbarred Galaxies

Our barred and nonbarred galaxy samples provide a unique testing ground to study the influence of the bar on the ISM. Figure 5 shows the distribution of the equivalent width of Hα at different radial bins for barred and nonbarred galaxies, revealing that barred galaxies tend to have less ionized gas than nonbarred galaxies in their inner regions. The ionization mechanisms of the gas in both samples are analyzed using BPT diagrams. Figure 6 shows barred and nonbarred galaxies on the [N II]/Hα BPT diagram with Wisconsin H-Alpha Mapper (WHAM) Milky Way observations overlaid (Haffner et al. 2003; Krishnarao et al. 2020). The MaNGA data are restricted to $S/N > 3$, which Belfiore et al. (2019) showed to be robust with minimal systematic errors for the emission lines of interest. To first order, barred galaxies show a larger spread in [N II]/Hα line ratios than their nonbarred counterparts. In the following sections, this relationship is further explored, and we consider radial trends in the ionized gas properties across different total stellar masses. Because only the [N II] BPT diagram is currently accessible to the Milky Way, our analysis from here on out is restricted to this diagnostic.

4.1. Radial Trends

Radial trends are considered for each galaxy sample after normalizing by their effective radii from the NSA catalog (Blanton et al. 2011) and correcting for inclination. Figure 7 shows the radial distributions of BPT classifications for all four samples of galaxies, revealing differences between the inner regions of barred and nonbarred galaxies. These violin plots show the probability density of the data smoothed by a Gaussian kernel and normalized across all classifications. The “Line-Less” category shows all spectra that have an Hα equivalent width $EW_{Hα} < 1$ Å and a g-band $S/N > 2$ isolating spaxels within the galaxy but with little to no gas content. Since MaNGA galaxies have more spaxels sampling the outer regions of galaxies than their inner regions, a violin plot of all spaxels in a MaNGA galaxy would have a wedge shape: narrow at small radii and wide at large radii. This style of figure shows qualitative differences in the distributions of classifications between samples. At any given radius, the relative widths of the distributions show the fraction of spaxels at those classifications.

Figure 3. Violin-style Gaussian kernel density plots of the stellar mass surface density at different radial bins of barred (orange) and nonbarred (blue) galaxies normalized by their effective radius for the analog (upper) and full MPL-8 (lower) samples. The interquartile range of each distribution is shown with a dashed line. The shapes of the violin plots are drawn based on the probability density distribution and have their widths normalized based on the number of spaxels in each category. The relative widths of the violin shapes across different categories at a given stellar mass surface density bin correspond to the fraction of spaxels within a category at that bin value. In general, both barred and nonbarred galaxies show the same distributions and trends.

correspond to the fraction of spaxels within a category at that bin value. While small differences can be seen in the inner $1R_e$ in general, both samples show similar radial distributions of their stellar mass.

Lastly, all of the barred and nonbarred galaxy samples are composed of galaxies from approximately 50% primary, 25% secondary, and 25% color-enhanced samples (see Wake et al. 2017), resulting in an equal sampling of the radial extent of galaxies. These tests confirm that differences seen between our barred and nonbarred galaxies should primarily result from the impact of the bar.

3.3. Milky Way Analogs

We apply a tight constraint on stellar mass based on the estimate of McMillan (2011) for the Milky Way of $M_*=6.43 \pm 0.63 \times 10^{10} M_\odot$. Systematic errors in the NSA stellar mass estimates are known to reach up to 0.2 dex (Pace et al. 2019b), so we loosen this constraint to encompass a range of $\pm 2.5\sigma$ of the McMillan (2011) estimate. This results in two “Milky Way analog” samples with and without a bar of sample size $N=62$ and 64, respectively.

The SDSS images of the 24 barred analogs within $1\sigma$ of the stellar mass criterion are shown in Figure 4. In this work, we will focus on the differences between the barred and nonbarred “Milky Way analog” samples and demonstrate the importance of a bar in defining analogs when considering the ISM. The larger full MPL-8 samples of barred and nonbarred galaxies will show the overall impact of a bar on the ISM across many galaxy scales. While more recent estimates of the total stellar mass of the Milky Way have shown lower values (e.g., Bovy & Rix 2013; McMillan 2017), we do not expect our selection of a specific mass range to significantly affect our results. All masses are considered throughout our analysis to help understand the importance of the total stellar mass in defining analogs.
In general, star formation is suppressed within the inner $1.5R_e$ of barred galaxies, where composite and L(N)ER classifications prevail. Figure 8 also shows this difference, where the barred galaxies have an increased presence at high [N II]/Hα and [O III]/Hβ line ratios within $R < 0.3R_e$. In the outer regions of galaxies, bars seem to have a lesser effect, and star formation dominates the radiation field. The L(N)ER classifications in nonbarred galaxies tend toward larger radii but do not show a strong difference from barred galaxies at these large radii. In our samples, AGN and Line-Less spaxels are rare and are not included in any further analysis.

### 4.2. Total Stellar Mass

Bars have a clear effect on the ionized gas in galaxies near the mass of the Milky Way and generally show the same effect across a wide range of stellar masses. However, this may be true for any stellar mass matched sample, so we compare the radial trends across mass bins in our full MPL-8 samples to consider the importance of the total stellar mass on this effect. Figure 9 shows an animation stepping through mass bins defined such that the sample size of barred galaxies remains constant at $N = 50$. Steps are made in increasing mass to add/remove three galaxies of greater/lower mass from the barred sample. The range in barred galaxy stellar masses in each step sets the galaxies used from the nonbarred sample. Similar to Section 4.1, radial distributions of BPT classifications and a Gaussian kernel density BPT diagram are shown to compare barred with nonbarred galaxies.

In general, galaxies across all masses show differences in their inner regions when considering the ionized gas. To test this claim, we use the Anderson–Darling statistical test with the null hypothesis that the fraction of BPT classifications as a function of radius are the same for both barred and nonbarred galaxies for all of our mass bins shown in the frames of Figure 9. Figure 10 shows the results of this test for three BPT classifications, with each pixel corresponding to the $p$-value of a test for a given radial $(x$-axis) and mass $(y$-axis) bin. Red points indicate where the null hypothesis can be rejected, indicating a statistical difference between barred and nonbarred galaxies. With the exception of the largest and smallest masses, the inner regions of barred and nonbarred galaxies within $R \lesssim 1.5R_e$ generally show a statistical difference in their distribution of BPT classifications, while outside of this range, the distributions are indistinguishable. The distribution of L(N)ERs within the range of Milky Way–like masses shows a statistical difference across a wider radial range than the other classifications. This may indicate that a loose cutoff mass exists to transition from galaxies in which the bar more significantly affects the ionization conditions of a galaxy. Similarly, star formation classifications show a more widespread difference at masses in the range of $10.4 \lesssim \log_{10}(M_*/M_\odot) \lesssim 10.8$. This mass range also corresponds well with signs of inside-out quenching as galaxies transition from star formation to quiescence, as shown in Belfiore et al. (2017; see their Figure 4).

![Figure 4](image-url)
To better understand how the presence of star formation differs between the inner parts of barred and nonbarred galaxies, we consider the fraction of spaxels within 1.5 \( R_e \) that have star formation BPT classifications in both samples across the same mass bins. The results of this are shown in Figure 11, where the barred galaxies show a decreased presence of star-forming spaxels across all masses. This difference is greatest in the range \( 10.5 \lesssim \log_{10}(M_* / M_\odot) \gtrsim 10.9 \), where the nonbarred galaxies show a 10%–30% greater presence of star-forming spaxels than barred galaxies.

### 5. Effect of Bar on Host Galaxy

Section 4 establishes an observable and significant impact on the ionized gas in galaxies with a bar compared to those without a bar. Honing in on the barred sample will allow for a more direct measure of how a bar impacts the ISM of its host galaxy. We compare the properties of the ISM interior and exterior to the bar radius and explore differences at a constant radius in the bar and interbar region. This is accomplished using bar masks from Galaxy Zoo:3D (see Section 2 for details).

Figure 12 shows differences in BPT classifications interior to versus outside of the bar radius across different masses for 226/240 barred galaxies in an animated plot. The gas interior to the bar radius shows greatly enhanced \([\text{N II}]/\text{H}\alpha\) line ratios than their nonbarred counterparts. Classification lines are from Kauffmann et al. (2003; orange solid line), Kewley et al. (2001; blue dashed line), and Schawinski et al. (2007; cyan dotted-dashed line).

### 6. Discussion

The LI(NE)R and, more generally, enhanced \([\text{N II}]/\text{H}\alpha\) emission line ratios in galaxies remain complicated and not...
fully explained since Burbidge & Burbidge (1965) noticed them in the inner regions of galaxies. The previous sections have shown an apparent link between the presence of LI(N)ERs and bars. This apparent link has also been observed in more nearby galaxy samples that find bar regions of galaxies to be largely LI(N)ERs (James & Percival 2015). The bar can be
interacting with and setting the conditions of gas in its vicinity in many ways, including increased shear (Renaud et al. 2013), shocks (Athanassoula et al. 1982), enhanced stellar densities (Anders et al. 2019), increased presence of evolved stars, and a nonaxisymmetric gravitational potential arising from the stellar structure.

Figure 9. Slice from an animated figure available in the online journal showing changes in the BPT classifications of barred (orange; top left) and nonbarred (blue; bottom left) galaxies across different stellar masses. Bins are defined such that the selected barred galaxy sample is fixed to 50 galaxies and stepped forward to add/remove three galaxies in each step. The nonbarred bins select galaxies within the min/max mass range of the barred galaxies, and their sample size varies. The pink and green shaded bars in the stellar mass histograms encompass the mass range used to define the current bin. Violin-style distributions of BPT classifications are shown as in Figure 7, and a Gaussian kernel density estimate BPT diagram for the binned sample is shown to the right for $R_e < 0.3$. This shows that the differences between gas in barred and nonbarred galaxies extend across a wide range in total stellar mass. A full description of this 13 s animation can be found in Appendix B.2. (An animation of this figure is available.)

Figure 10. Results of the Anderson–Darling statistical test that the underlying distribution of the fraction of spaxels in different BPT classifications comes from the same parent distribution across different stellar mass and $0.2 \, R_e$ radial bins. The transition from red to blue pixels is set at the $p = 0.05$ level; deep red points signify that the null hypothesis can be rejected at the $p = 0.01$ significance level and indicate a statistical difference in the fractional classifications. The four panels correspond to the BPT classifications of star formation (first panel), composite (second panel), and LI(N)ER (third panel). The AGN classifications are not considered here, since they are an overall rare occurrence in these samples. Stellar mass bins are defined as in Figure 9 such that each mass bin has a fixed barred sample size of 50 barred galaxies, with larger masses situated higher along the $y$-axis. The median stellar mass and the range of masses considered in each bin are plotted in the fourth panel. The dashed yellow line encloses samples that include a Milky Way stellar mass value of $M = 6.43 \times 10^{10} \, M_\odot$ (McMillan 2011). We interpret these results to show that the inner regions of barred galaxies, especially around $0.2–1.5 \, R_e$, have a statistically significant difference from nonbarred galaxies in their ionized gas properties across all but the highest and lowest stellar masses.
Inflowing material along the bar can fuel central mass concentrations (such as the CMZ in the Milky Way) or AGN, as observed in other galaxies (e.g., George et al. 2019) and in simulations (e.g., Athanassoula 1992). The role of bars in fueling AGN has been largely debated, as their role may not be vital (Heckman & Best 2014) and there are only slight enhancements of AGN in barred galaxies compared with nonbarred galaxies (Galloway et al. 2015). The increased presence of Li(N)ER-like spectra in bars can complicate such studies, as AGN may be wrongly identified if using limited diagnostic methods.

A proposed source of ionization for Li(N)ERs, and gas with enhanced line ratios in general, is the classes of hot evolved stars, such as post–asymptotic giant branch (AGB) stars, subdwarf OB (sdOB) stars, or K and M giants (Burbidge & Burbidge 1965; Yan & Blanton 2012; James & Percival 2015; Belfiore et al. 2016). However, these sources may not emit enough Lyman continuum radiation to account for the ionized gas present, as shown locally in the Milky Way (Reynolds 1989), or fail to predict all of the observed line ratios (Yan 2018). The K and M giants as a source of ionization, as described in Burbidge et al. (1963), is distinct in that they provide a source of heating through ejected winds, as opposed to purely photoionizing the gas. This additional heating can help explain the large [N II]/Hα line ratios as an increase in gas temperature (up to 20,000 K). This mechanism can be effective in the central parts of galaxies, where a large density of these older stars may be present, but would likely not make a significant impact in outer Li(N)ERs (e.g., James & Percival 2015; Belfiore et al. 2016), where photoionization may again dominate.

If hot evolved stars were primarily responsible for the observed ionization, a high density of stars would be required to produce enough ionizing photons from these relatively low-luminosity sources. While at times (e.g., within the bar) the presence of Li(N)ERs is well correlated with large stellar mass surface densities, there are still some regions outside of a bar or generally at areas of low stellar mass surface density that are Li(N)ERs. It may still be possible for a centralized or other distant and luminous source to ionize the gas (Domgorgen & Mathis 1994; Wood & Mathis 2004), especially in a region of the galaxy where the ISM may be evacuated by the bar (Sormani & Barnes 2019), providing channels for photons to travel large distances. More in-depth modeling of the evolution, radiation field, and distribution of sdOB stars (see Heber 2016, for a review) is necessary to fully constrain the possible ionization sources of Li(N)ERs.

Bars show a strong correlation with the physical conditions of the ionized ISM, with bars generally affecting the gas state within the bar radius and, more specifically, gas in the bar and interbar region (see Figure 12). The use of Galaxy Zoo:3D bar masks reveals these differences and demonstrates a clear nonaxisymmetric impact of the bar to the ionized gas of galaxies. The presence of Li(N)ER gas in the interbar region is especially difficult to ionize from a population of evolved stars, since the stellar surface mass density in these regions is generally lower. However, this gas may still feel the effects of an increased population of hot evolved stars within the bar if radiative transfer is able to propagate enough photons radially outside of the bar. In-depth 3D Monte Carlo radiative transfer models (see, e.g., Wood et al. 2010; Vandenbroucke et al. 2018) may provide better constraints on the ionization mechanisms at play in the vicinity of a bar. Shocks may also be a viable source of ionization for Li(N)ERs, especially since dust lanes within bars are likely formed via shocks (Athanassoula 1992; Lindblad et al. 1996; Lindblad & Kristen 1996; Englimaier & Gerhard 1997; Greve et al. 1999).

Better diagnosing of the presence of shocks may be possible with use of the [S II] λ6716 emission line in combination with the [N II] and Hα lines and can soon also be tested in the Milky Way using upcoming multiwavelength WHAM observations of the tilted disk (see Krishnarao et al. 2020). Higher spatial resolution IFU studies of face-on systems, such as the TIMER (Gadotti et al. 2019), PHANGS (Sun et al. 2018; Utomo et al. 2018), or TYPHOON (Poetrodjojo et al. 2019) surveys, may be able to better identify these shock fronts and test for their contributions toward the overall ionization and physical conditions. These nearby galaxies also provide a venue to study the resolved behavior of neutral atomic and molecular gas in H I and CO. While large-scale programs to characterize the total H I content of MaNGA galaxies exist (HI-MANGA; Masters et al. 2019), only limited and targeted efforts using interferometric observations look to provide spatially resolved information (e.g., Newnham et al. 2019, for non-MaNGA galaxies). Understanding the possible role of a galaxy’s environment on both bars and their effects on the ISM using existing value-added catalogs, such as the Galaxy Environment for MaNGA Galaxies (GEMA) catalog (e.g., Argudo-Fernández et al. 2015), is also possible. However, MaNGA does not sample many galaxies in dense clusters, where these environmental effects would be best studied, so such an analysis may be more suitable for other galaxy surveys.

Many of the predictions from simulations (e.g., Athanassoula 2002; Renaud et al. 2013; Seo et al. 2019) and previous observational efforts (e.g., Athanassoula et al. 1982; Masters et al. 2011, 2012) are consistent with the effects seen in this work, such as suppression of star formation around a bar, with star formation arising in rings at larger radii, or in central cores, like in the CMZ (Longmore et al. 2013). The Milky Way analogs tend to show consistent behavior with that of the Milky Way, in particular the presence of Li(N)ER classifications within the bar and star formation dominating the radiation field outside the bar. Combining the predictions from simulations,
observations in the Milky Way, and analysis of Milky Way analogs in this work, we can speculate on the face-on distribution of ionized gas and dominant ionization mechanisms in the inner Galaxy within the bar radius.

Figure 13 shows a predicted cartoon schematic of how a barred galaxy like the Milky Way may appear when considering the ionized gas conditions. Gas can be funneled inward toward the center, where it can condense and form stars in a CMZ-like core. Outside of this region, the gas has large [N II]/Hα line ratios and is largely Li(N)ER, especially within the nonaxisymmetric structure of the bar. Along the bar’s minor axis but outside of the bar itself, gas transitions from Li(N)ER to composite classifications as the [N II]/Hα line ratio decreases. At around the bar radius, star formation dominates the radiation field in a ringlike form and persists throughout the rest of the galaxy as the primary source of ionization.

7. Conclusions

Especially in early-type (Ho et al. 1997) and spiral galaxies, Li(N)ERs have long been seen to be a common occurrence, with early estimates showing up to 1/3 of spiral galaxies are Li(N)ERs (Heckman 1980). Here we show that among disk galaxies, barred galaxies show a much greater occurrence of central Li(N)ER-like spectra when compared with nonbarred galaxies across all but the smallest and largest masses. This demonstrates an apparent link between the presence of a bar and conditions of ionized gas associated with Li(N)ER-like spectra. Many mechanisms can be at play in the environment of a bar, ranging from large shearing motions
and shocks commonly observed to suppress star formation. Burbidge & Burbidge (1965) first suggested that gas showing such large ratios of $\text{[N II]}/\text{H}\alpha$ may originate from gas that fails to condense, as would be the case for the inner Milky Way and other barred galaxies where the star formation efficiency is low (Longmore et al. 2013; Barnes et al. 2017).

Regardless of the specific mechanisms at play, it is clear that bars have a significant impact on the ionized gas and radiation field in galaxies. Barred galaxies show at least a 10% decrease in the presence of star formation BPT classifications in their inner $1.5R_e$ when compared with nonbarred galaxies for all but the highest stellar masses. At around the stellar mass of the Milky Way, this effect is stronger, with a difference of 20%–30%. The presence of a bar may be one of the most important criteria for defining a sample of Milky Way analogs when investigating the ISM, especially in inner regions of galaxies. Additionally, the total stellar mass of a galaxy seems to have relatively little effect on the state of the ionized gas for intermediate masses around that of the Milky Way. This suggests that the total stellar mass is not a very important criterion in defining our analog sample. Studying a large sample of face-on barred galaxies has allowed us to better understand and predict the behavior of the ISM in the Milky Way. Continued studies of the dynamics and physical conditions of gas in barred galaxies and the Milky Way, combined with observations of optical emission and ultraviolet absorption and submillimeter-to-radio wavelength studies of the inner Milky Way can provide the highest-resolution counterpart to studies of face-on analog samples.

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Facilities: GALEX, Sloan, WHAM.

Software: astropy (Robitaille et al. 2013; Astropy Collaboration et al. 2018), matplotlib (Hunter 2007), numpy (van der Walt et al. 2011), scipy (Jones et al. 2001), seaborn (Waskom et al. 2014), sdss-marvin (Cherinka et al. 2019), pandas (McKinney 2010).

**Appendix A**

**Galaxy Zoo:3D—Bar Length Validation**

Bar lengths measured using Galaxy Zoo:3D bar masks are validated with an independent measure of the bar length. Bar lengths are measured following the methodology employed in Kraljic et al. (2012). This method decomposes stellar surface density profiles inferred from white-light images into

![All Bar Length Methods](image)

**Figure A1.** Histogram of bar length measurements from the Fourier method employed in Kraljic et al. (2012; pink) and from bar masks using Galaxy Zoo:3D with vote thresholds of 10% (green), 20% (orange), 30% (blue), and 40% (yellow) in the left panel. The fraction of galaxies from the barred MPL-8 sample that have valid bar masks with each threshold is shown in parentheses. Increasing the vote threshold parameter increases the agreement between bar length measurements from both methods but greatly decreases the number of galaxies that return valid bar masks. The right panel shows histograms of the differences in bar length measurements between the Fourier method and Galaxy Zoo:3D, color-coded by vote thresholds.
azimuthal Fourier components. The signature of a bar is seen in the even \((m = 2)\) Fourier component within the bar region. Bars can be identified when the \(\Phi_2(r)\) phase is constant with radius. This method requires a reliable estimate of the bulge-to-total ratio (BTR) of the host galaxy to correctly guess the bar starting position. The BTR \(r\)-band measurements from two-component decompositions of Simard et al. (2011) are used for these initial guesses. The presence of an additional bar component can bias the amount of light attributed to the bulge and hence give a BTR that is too large for a galaxy. When this happens, the starting position of the bar-finding code is too large in radius, and the bar may not be correctly identified. Out of the 240 barred galaxies in our full MPL-8 sample (described in Section 3), 206 returned reasonable measurements.

Figure A1 shows a histogram of the measured bar lengths using the Fourier method and various vote thresholds for Galaxy Zoo:3D. The vote threshold parameter is the percent of votes that must be present to flag a spaxel as part of the bar mask. Increasing this threshold selects fewer spaxels in a galaxy but can help to ensure the bar mask is isolated to spaxels with a high confidence of being in the bar. Increasing this threshold also reduces the total number of galaxies that have valid bar masks. In the right-hand panel of Figure A1, the difference between the Fourier bar length measurement and the bar mask measurements for different vote thresholds are shown. Increasing the threshold improves the agreement between both methods, suggesting that too many spaxels are selected for the bar mask with low thresholds. However, the completeness of selecting barred galaxies greatly diminishes with larger threshold values. This is seen by the fraction of galaxies that return a valid bar mask for each threshold, shown in square brackets in the Figure A1 legend. To balance these effects, we use a vote threshold of 20% in this work. Previous work using earlier versions of Galaxy Zoo:3D also used a 20% vote threshold (e.g., Fraser-McKelvie et al. 2019).

Appendix B

Animated Figure Descriptions

Full descriptions of the contents of the animated figures used in this work are available below.

B.1. Figure 8

The animation in Figure 8 runs for 10 s and steps through radial bins starting at 0.0 < \(R/R_e\) < 0.3 and ending at 2.9 < \(R/R_e\) < 3.1. The distributions of both barred and nonbarred galaxies in both the analog and full MPL-8 samples have the largest dispersion in the innermost radial bin (corresponding to the still image in the print version). At larger radii, the presence of LI(N)ER classifications decreases, star-forming classifications increase, and in general, the [N II]/H\(\alpha\) line ratio decreases. At large radii, both distributions move toward the upper left corner of the diagram within the star-forming and composite classifications. The greatest difference between the barred and nonbarred samples is seen within \(\sim 1 R_e\), with barred galaxies showing a significantly larger width along both axes, compared to a more strongly peaked distribution for nonbarred galaxies.

B.2. Figure 9

The animation in Figure 9 runs for 13 s and steps through total stellar mass bins starting at 9.29 < \(\log_{10}(M/M_\odot)\) < 10.13 and ending at 10.99 < \(\log_{10}(M/M_\odot)\) < 11.35. As the animation progresses, the stellar mass distributions in the left panel remain the same, and the pink and orange bars slide from left to right to encompass the current range of masses being considered in the middle and right panels. The violin distributions in the middle panel show that star-forming classifications are equally dominant in the inner parts of barred and nonbarred galaxies at the lowest masses, while their presence in barred galaxies decreases at larger stellar masses. Composite classifications are uncommon in the lowest masses and gradually increase in their presence at larger masses. The LI(N)ER classifications show a larger presence in barred galaxies, especially within \(1.5 R_e\) at lower masses, and continue to show this trend at almost all mass bins but the very largest. In the right panel, the BPT diagram is concentrated primarily on the star-forming classification at the lowest masses and expands to encompass composites and LI(N)ERs at larger masses.

B.3. Figure 12

The animation in Figure 12 runs for 13 s and steps through total stellar mass bins as in Figure 9. The distribution of total stellar mass in the top left panel remains the same, and the pink bar slides from left to right to encompass the current range of masses being considered in the other panels. In the bottom left panel, star-forming classifications initially dominate at the lowest masses and decrease at larger masses, becoming relatively rare in both the bar and interbar regions. Composite classifications start very low and increase to dominate the interbar region at all radii at larger masses. The LI(N)ER classifications also start very low and then rise at larger masses to dominate the inner regions of both the bar and interbar region. The middle panel shows the same overall trends but now considering all spaxels equally within and outside the bar radius. At the lowest masses, the animation begins to show that the star-forming classification dominates at all radii and gradually decreases, especially within the bar radius, at larger masses. Composite classifications gradually rise in prominence, especially near and beyond the bar radius. The LI(N)ER classifications rise to dominate within the bar radius and are also somewhat present at larger radii at larger masses. The right panel shows that the bar, interbar region, and region outside the bar radius are relatively similar at the lowest masses. As the animation progresses, the region outside the bar radius remains relatively in place, while the bar and interbar distributions spread out along both axes. The bar region undergoes this spread faster and develops a peak in its distribution at larger values of the line ratios.

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