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Light-induced current oscillations in the charge-ordered state of (TMTTF)$_2$SbF$_6$

Tobias Peterseim and Martin Dressel*

1. Physikalisches Institut, Universität Stuttgart, Pfaffenwaldring 57, D-70550 Stuttgart, Germany
* dressel@pi1.physik.uni-stuttgart.de, phone +49 711 6856 4946 and Fax +49 711 6856 4886.

Abstract: Below $T_{CO} = 157$ K the quasi-one-dimensional charge-transfer salt (TMTTF)$_2$SbF$_6$ undergoes a pronounced phase transition to a charge-ordered ground state. We have explored the non-linear and photoconductive behavior as a function of applied voltage, laser pulse energy and temperature. Besides a decay of the photoconductive signal in a double exponential fashion in the millisecond range, we discover current oscillations in the kHz range induced by the application of short laser pulses. While the resonance frequencies do not depend on voltage or laser intensity and vary only slightly with temperature, the amplitude changes linearly with the laser intensity and voltage. The findings are discussed and compared to comparable phenomena in other low-dimensional electron systems.

Keywords: charge-transfer salts; TMTTF; charge-order phase transition; non-linear transport; photoconductivity; current oscillations; calculations of band structure and optical conductivity

1. Introduction

The charge-transfer salts (TMTTF)$_2$X, where TMTTF stands for tetramethyltetrathiafulvalene and X for a monovalent anion, have been under scrutiny due to their quasi-one-dimensional electronic character and ordering phenomena in the spin and charge sectors [1–4]. During the last years several comprehensive investigations have been conducted on the optical [5–9] and transport properties [10], the lattice [11,12] and electronic structure [13–15] as well as on the magnetic properties [16–21]. However, not much is known about the non-linear properties at the charge-order transition taking place at $T_{CO} = 157$ K and transient conduction phenomena. This is surprising, since nonlinear behavior was observed in various inorganic as well as organic compounds in the vicinity of a metal-insulator transition, such as the charge-density wave transition in NbSe$_3$, the metal insulator transition in VO$_2$, or the charge-order transition in the two-dimensional $\alpha$-(BEDT-TTF)$_2$I$_3$ salt [22–28].

Several classes of electrically induced nonlinear phenomena in solid-state materials can be distinguished:

- Non-equilibrium charge carriers: modification of the charge carrier number or the mobility.
- Space-charge distribution: for instance impurity bands.
- Electrically induced phase transitions: modification of the electronic system or the crystalline structure.
- Ionic motion in the host material: forming of conductive filaments.
- Sliding charge density waves or charge order.
- Domain wall motions or excitations of solitons.

Although exotic electrically induced effects emerge in most classes of materials, it is basically impossible to draw a conclusive picture since the diversity of the involved ground states is enormous. Nevertheless, in regard to future electronic devices and applications, more effort has to be devoted to the experimental studies on electrically induced phenomena as well as the development of accurate models explaining a wide range of these effects occurring in various material classes. For a broader discussion we refer to the review articles [29,30] and textbooks [31,32] and further literature listed therein.
Figure 1. Experimental characterization of (TMTTF)$_2$X, X=PF$_6$ (black), AsF$_6$ (red), and SbF$_6$ (green) single crystals. (a) Temperature-dependence of the resistivity along the $a$-direction [10]. The labels mark the charge-order transition temperature $T_{\text{CO}} = 67, 102, \text{and} 157 \text{ K}$, respectively. The inset demonstrates the effect of hydrostatic pressure on the transition by displaying the temperature-dependent dc resistivity of (TMTTF)$_2$SbF$_6$ measured along the $b$-direction. $T_{\text{CO}}$ is seen as a kink in $\rho(T)$ that shifts to low temperatures as pressure increases. (b) Temperature-dependence of the spin susceptibility at constant volume $\chi_s$. (TMTTF)$_2$PF$_6$ and (TMTTF)$_2$AsF$_6$ undergo a spin-Peierls transition at $T_{\text{SP}} = 19 \text{ K}$ and $13 \text{ K}$, respectively. (TMTTF)$_2$SbF$_6$ exhibits an antiferromagnetic state below $8 \text{ K}$ causing an increase of $\chi_s$. Above the transition $\chi_s$ can be described by a model of a 1/2 antiferromagnetic Heisenberg chain [16,38]. (c) The real part of the dielectric constant $\epsilon_1$ measured at $f = 100 \text{ kHz}$ as a function of temperature $T$ reveals a Curie-law behavior [39]. (d) $^{13}$C-NMR spectra of (TMTTF)$_2$AsF$_6$ at $200 \text{ K}$ and $80 \text{ K}$. Above $T_{\text{CO}} = 102 \text{ K}$ two distinct features A and B can be identified due to the two molecules per unit cell. They split into a total of four peaks, as depicted by the arrows [40], giving evidence for charge order.

The first studies on electrically induced phenomena in organic conductors were performed in the 1970's and 1980's on Cu(TCNQ) [33] and TTF-CA [34]. Later, this was extended to other organic systems, such as K(TCNQ) [35], where the switching was also recorded optically, expressing a stripe structure of alternating insulating and metallic domains between the two contacts. In the two-dimensional organic salt $\theta$-(BEDT-TTF)$_2$CsZn(SCN)$_4$ [36] and its sister compounds [37] the charge order was melted by applying a current to the crystal [30]. To date, no concept or model could be established that can explain all these observations in organic salts.

Since nonlinear behavior was observed in several different inorganic as well as organic compounds at the vicinity of a metal-insulator transition, we expect similar phenomena in the quasi one-dimensional salts (TMTTF)$_2$X, with the centrosymmetric anion X=PF$_6$, AsF$_6$, and SbF$_6$. In Figure 1 we collect the important experimental findings characterizing these three materials. Upon cooling the resistivity goes through a shallow minimum due to electronic localization and later exhibits a kink at the charge-order transition $T_{\text{CO}}$, as displayed in panel (a). No significant change is observed in the magnetic susceptibility at $T_{\text{CO}}$ [17] plotted in Figure 1(b); while the magnetic order...
Figure 2. Temperature dependence of the charge disproportionation probed by the splitting of the infrared-active $\nu_{28}$ molecular vibrational mode in (TMTTF)$_2X$ with $X = \text{PF}_6^-$, $\text{AsF}_6^-$, $\text{SbF}_6^-$, and $\text{TaF}_6^-$ as a function of temperature $T$ recorded along the $c$ axis. The charge imbalance $2\delta$ is estimated from the difference of the two resonance frequencies by $2\delta = \Delta \nu / (80 \text{ cm}^{-1}/e)$ [7,41].

Figure 3. Two views on the crystal structure of (TMTTF)$_2\text{SbF}_6$ illustrate the confinement of the octahedral anions in a cavity formed by methyl groups (gray) of the surrounding TMTTF molecules. Also shown are the closest connections between the fluorine (green) and sulphur atoms (yellow).

and spin-Peierls transitions occur at low temperatures. Due to charge order, the dielectric constant exhibits a maximum at this temperature (see panel (c)), more or less following Curie’s law [39]. As presented in Figure 2, the charge-sensitive vibrational modes split due to charge disproportionation [7]; similar conclusions are drawn from the NMR spectra [40] plotted in Figure 1(d).

For (TMTTF)$_2\text{SbF}_6$ a surprising behavior of the material was observed in response to electrical and optical pulses, which is presented and discussed below in detail. We could not detect similar phenomena in (TMTTF)$_2\text{PF}_6$ and (TMTTF)$_2\text{AsF}_6$. As depicted in Figure 1(a) and (b), (TMTTF)$_2\text{SbF}_6$ reveals the steepest slope at the phase transition, while (TMTTF)$_2\text{PF}_6$ and (TMTTF)$_2\text{AsF}_6$ exhibit a smooth transition. Therefore, we conclude that for the one-dimensional salts strong resistivity changes [10] at the transition temperature are a good indication that nonlinear behavior can be observed. Otherwise, the difference is too small and the photo-induced effect is absent.

2. Experimental Details

Single crystals of the charge-transfer salts (TMTTF)$_2\text{SbF}_6$ were grown by electrochemical methods as described previously [10,38]. The needle-shaped single crystals are several millimeters long in the $a$-direction and less than a millimeter wide in the other crystal directions. Figure 3 shows the crystal structure of (TMTTF)$_2\text{SbF}_6$ in different orientations.
Figure 4. Schematic drawing of the photoconductivity setup. It consists of a pulse generator creating the voltage $U_S$ at the sample, an oscilloscope, and a load resistor. $R_L$ is used to determine the current in the circuit. The frequency doubled Nd:YAG laser beam is focussed through the cryostat window.

For transport experiments four fine gold wires were attached to the crystal by carbon paste as contacts. The samples were cooled in a helium cryostat. Besides simple $\rho(T)$ characterization (Figure 5), we measured the transport response of the crystals when a strong electric pulse is applied.

In a second step, photoconductivity was studied, i.e. we probed the time response of the electric transport after exposing the crystal to a laser pulse (8 ns) with a wavelength of $\lambda = 532$ nm (corresponding to a photon energy of 2.35 eV). The setup for the photocurrent measurements sketched in Figure 4 consists of two contacts on the specimen separated by 1.1 mm. A voltage pulse with variable width is created on the contacts by a pulse generator. The applied voltage $U_S$ and current $U_S/R_L$ are measured by an oscilloscope. Since the laser has a low repetition rate between 1 Hz and 20 Hz heat accumulation is avoided; it is also ensured that long living states decay before the next pulse arrives and hence, accumulation effects can be excluded. The generated photo current $I_{ph}(t)$ can be derived from the following physical relation. The electronic circuit consists of a voltage source, the sample and the load resistance $R_L$. The dark current $I_D(t)$ is defined as

$$ I_D(t) = \frac{U_S}{(R_L + R_{S,D})} \quad ; $$

$R_{S,D}$ the dark resistance of the sample and $R_{S,Ill}$ the corresponding resistance when illuminated. The total current under illumination $I_{Ill}(t)$ is respectively

$$ I_{Ill}(t) = \frac{U_S}{(R_L + R_{S,Ill})} \quad .$$

with $R_{S,Ill}$ Thus, the measured photo current is

$$ I_{ph}(t) = I_{Ill}(t) - I_D(t) = U_S \left( \frac{1}{(R_L + R_{S,Ill})} - \frac{1}{(R_L + R_{S,D})} \right) \quad .$$

In a photocurrent measurement the contact resistance plays a crucial role since it influences the transport of the charge carriers to the contact. The materials, i.e. metal, semiconductor or insulator, touching at the interface form different types of contacts. At the boundary the band structure is deformed due to the difference in work function and electron affinity of the materials. Therefore, a barrier is created at the boundary surface leading to a depletion area [31,42]. By that, additional carriers can be transferred or the charge carriers can be blocked at the contacts. There are typically two kinds of contacts: the Schottky-contact and the Ohmic contact. Generally, several processes contribute to the charge transport across the barriers. The dominating processes are the thermal emission of charge carriers across the energy barrier and tunneling of the carriers, being temperature-independent. However, it is assumed that Ohmic contacts are formed in the organic crystals. During the transport measurements the temperature-dependence of the contacts
was checked by measuring the contact resistance at specific intervals. Their resistance value normally changes at a phase transition when a gap opens in the electronic system. For details on the time and field-dependent experiments we refer to Refs. [43,44].

3. Response to Electric Fields

Before discussing the response of (TMTTF)$_2$SbF$_6$ to electric pulses, let us consider the temperature-dependent resistance from room temperature down to 125 K, in order to determine the phase transition temperature. This allows us to assign $T_{CO}$ more precisely compared to optical methods [7,8] and to deduce the contact resistances in this particular case. In Figure 5 the logarithm of the resistance of (TMTTF)$_2$SbF$_6$ is illustrated for the $a$-direction in a temperature range from room temperature down to 125 K. When lowering the temperature, the resistance decreases down to a minimum at $T = 215$ K, which corresponds to the sign change of its slope (see inset). This is caused by electron-electron interactions opening up a Mott gap at the middle of the upper half-filled band and Umklapp scattering. Thus, the resistance further increases with decreasing the temperature down to the charge-order transition, corresponding to a linear slope of the derivative. At $T_{CO}$ a kink arises and below this temperature $\ln R$ increases nonlinearly; hence transport is not just thermally activated. It is worth to mention that for (TMTTF)$_2$SbF$_6$ the resistivity increases more abruptly and steeper at the charge-ordered phase transition than observed in the (TMTTF)$_2$PF$_6$ and (TMTTF)$_2$AsF$_6$ analogues. Furthermore, (TMTTF)$_2$SbF$_6$ reveals an anomaly in the thermal expansion coefficient at the charge-order transition [45]. In contrast, in its sister compounds no feature appears at the corresponding transition temperature.

![Figure 5](image-url)  
**Figure 5.** Logarithm of the resistance of (TMTTF)$_2$SbF$_6$ as a function of the inverse temperature $1/T$. At about $T = 215$ K a minimum appears, marking the transition to the localized state related to the opening of a Mott gap. The resistance increased slightly and at $T_{CO} = 157$ K a kink is observed in the curve due to the charge order transition. Inset: derivative of the logarithm of the resistance vs. $1/T$. The charge localization leads to a change of the sign from negative (metallic) to positive (insulating). Also the slope changes at $T_{CO}$ leading to a rather flat curve at low temperatures.

A similar behavior appears in the temperature-dependent evolution of the charge disproportion [7], as depicted in Figure 2. While in (TMTTF)$_2$PF$_6$ and (TMTTF)$_2$AsF$_6$ a mean-field behavior appears, a first-order like transition takes place in (TMTTF)$_2$SbF$_6$. The latter is probably related to the enhanced dimerization (see [10,12]) and the charge localization state at elevated temperatures. This behavior is even more pronounced in (TMTTF)$_2$TaF$_6$, where the phase transition takes place at $T_{CO} = 177$ K and $2\delta = 0.28e^\dagger$ is of comparable size [41].

We now apply voltage pulses of 1 ms duration to the sample and measure the total current and the voltage drop across a load resistor of $R_L = 1$ k$\Omega$ connected in series to the sample. In Figure 6...
the time-dependent current response of (TMTTF)$_2$SbF$_6$ at $T = 140$ K is plotted for two different voltages. At the beginning of the pulse, an overshoot of the measured current occurs due to the parasitic capacity and resistance in the circuit (panel a). The short oscillating signal located at 100 µs is caused by a parasitic current induced by the switching of the Pockels cell of the laser system, later used to create the photocurrent. For low voltages ($U = 0.9$ V) the current signal stays constant during the whole pulse. When increasing the amplitude of the pulse, the current increases approximately in a linear way obeying Ohm’s law. Nevertheless, following the temporal evolution further reveals a nonlinear increase of the current, which can be attributed to a thermal modification of the contact resistance.

Figure 6(b) displays current pulses for different temperatures below $T_{CO}$. The current signal is reduced for the same applied voltage in the same way as the total resistance increases with temperature. Independent on the applied voltage or examined temperature no switching behavior or current oscillation can be detected that would indicate a sliding of the $4k_F$ charge-density wave (CDW) or a collective response of the charge-ordered state [46,47].

4. Response to Laser Pulses

Next the photo response of (TMTTF)$_2$SbF$_6$ was studied by illuminating a fraction of the space between the two contacts by a laser pulse. In Figure 7 the recorded current pulse is illustrated exemplarily at $T = 150$ K for a 5 V voltage pulse and a laser pulse with the energy of 40 µJ. Similar to Figure 6, an overshoot is present at the beginning. Additionally, 750 µs after the onset of the voltage pulse, the current suddenly shoots up with a rise time a few nanoseconds. Strikingly, this onset is correlated with the incident laser pulse, giving clear evidence for photoconductivity. The signal barely decays and stays almost constant over the entire voltage pulse length. In the inset of Figure 7 an enlarged detail of the first 400 µs of the photocurrent is visualized. Surprisingly, current oscillations are superimposed on top of the signal, whose amplitude decays with time. The inset further reveals a slight signal decay within the first 59 µs.

We examined the photocurrent signal for various voltages below the phase transition $T_{CO} = 157$ K, as presented in Figure 8(a) for the temperature $T = 150$ K and laser pulse energy $E = 32$ µJ. Above $T_{CO}$, no change was observed. There, the photocurrent was extracted by subtracting the
Figure 7. Current response to a 1 ms long 5 V voltage pulse applied on a (TMTTF)$_2$SbF$_6$ crystal at a temperature of 150 K. At about $t = 0.075$ ms, an additional laser pulse of $E = 40 \, \mu$J (black arrow) illuminates the sample. This induces a photocurrent, visible in the steep, weak jump of the current signal. In the inset the first 400 $\mu$s of the photocurrent response are magnified. Interestingly, overall signal is are superimposed with pronounced oscillations.

Figure 8. Time-dependent photocurrent of (TMTTF)$_2$SbF$_6$ for (a) various voltages, (b) laser pulse intensities and (c) temperatures as indicated. In all three panels the temporal dynamic is fitted by a double exponential function (black lines).
current without any light (see for example Figure 6) from the total current signal. Thereby we can also eliminate the nonlinear transport behavior of the device, consisting of the contact resistance and the sample. By enlarging the applied voltage the signal increases linearly. Similarly, the photocurrent becomes larger with rising laser pulse energy and the temporal behavior is the same as for the voltage-dependent current curves, as depicted in Figure 8(b). Finally, also the temperature dependence of the transient photocurrent is plotted for a fixed pulse energy 32 µJ and voltage amplitude $U = 5$ V in Figure 8(c). The photocurrent is the smallest in the vicinity of the phase transition ($T = 154$ K) and then stays almost constant below $T = 150$ K. One should note that all curves in the latter three graphs exhibit the same temporal dynamics, which can be fitted by a double-exponential decay function $f(t) = y_0 + A \exp \{t/\tau_1\} + B \exp \{t/\tau_2\}$. Between $T = 154$ K and 125 K, the first time constant $\tau_1$ decreases from 30 µs to 16 µs and $\tau_2$ diminishes by a factor of 2 from 300 µs to 170 µs. Here $y_0$ is the background which can be considered as a third extended living state, decaying at least on a millisecond time range. Such relaxation dynamics have also been observed in other photo-response experiments on organic conductors [48–51] taking place between the microsecond regime and far above the millisecond time scale.

However, the most important phenomena are the interference patterns which are superimposed on top of the exponentially decaying current signal. To the best of our knowledge, this has never been detected before in a photoconductivity experiment on organic conductors. In order to analyze the oscillations in detail, the frequency components were extracted from the decaying current by fast Fourier transformation (FFT). The amplitudes of the frequency components from the curves, depicted in Figure 8, are visualized in Figure 9 for various parameters. Five resonance frequencies can be identified: 49 kHz (labeled A), 54 kHz (B), 58.5 kHz (C), 77.5 kHz (D), and 82 kHz (E). Varying the voltage does not influence the resonance frequency (panel a), although the amplitude increases linearly. The same is true for the laser intensity dependence of the oscillation amplitude; it increases linearly without any modification of the peak positions, as shown in Figure 9(b) for $T = 150$ K and $U = 5$ V. It is important to note that the dark current (0 µJ laser power) is independent of the applied voltage, as it is exemplarily displayed in Figure 6; no oscillations occur, resulting in a flat amplitude
spectrum. While the resonances display only an amplitude-dependence on the voltage and pulse intensity, both the amplitude as well as the resonance frequency depend on temperature.

5. Discussion

We now turn to the question of the origin of the photocurrent and its dynamics: Why does the current oscillate after illumination?

Iwase et al. proposed a phenomenological description \[ \sigma_{\text{dc}}(T,J) = \sigma_1(T) + \sigma_2 J^n \]

where the first term on the right side describes the thermal activated linear Ohmic behavior and the second term accounts for the nonlinear conductivity. A two-state model assumes that “hot electrons” are generated in the conduction band similar as in inorganic semiconductors. Mori and collaborators \[54,55\] successfully applied this approach to \( \beta''-(\text{BEDT-TTF})_3(\text{HSO}_4)_2 \) and \((\text{TMET-TTP})_4\text{PF}_6\). The electron system is assumed to be decoupled from the lattice; hence an effective electron temperature \( T_e \) can be defined. Note, this is just a hypothetical temperature and does not mean that the sample is heated up. These excited electrons or quasi-particles follow the heat balance equation:

\[
n C_p \frac{dT_e}{dt} = P + \nabla(\lambda_{\text{therm}} \nabla T) ,
\]

which can be simplified by assuming a linear spatial heat transfer, expressed by the last term, leading to:

\[
n C_p \frac{dT_e}{dt} = P(t) - \lambda_{\text{therm}}(T_e(t) - T_0) .
\]

The left side of the equation is the energy stored in the excited electronic system. The first term on the right hand side is the electric input power \( P \) in W cm\(^{-3}\) created by the externally applied electric field or current, which can be expressed as follows:

\[
P(t) = E \times J(t) = \sigma(T_e(t))E(t)^2 \quad \text{or} \quad = \frac{J(t)^2}{\sigma(T_e(t))} ,
\]

and the second term of equation 6 describes the linear heat transfer from the electronic system to the lattice, which is kept at the initial temperature \( T_0 \). Here \( \lambda_{\text{therm}} \) is the thermal conductivity and \( nC_p \) the electron heat capacity per volume. Equation 6 is solved numerically for a certain starting electric field or current pulse. In general \( \sigma(T_e) \) is determined from the conductivity curve of the corresponding material because the excited carriers are responsible for the dc conductivity. From this model, a delay time \( \tau_d \) can be estimated which also accounts for current oscillations.

The possible influence of thermal effects on the induced current has to be discussed. Following the approach presented in Ref. \[28,56\], we calculated the temporal behavior of the sample temperature and correlate it with the temperature-dependent sample resistance depicted in Figure 5. The result of the simulation is visualized in Figure 10 for various temperatures \( T \) between 125 K and 154 K. We assume that the complete pulse energy is absorbed and transferred into Joule heat at the surface of the specimen as the penetration depth is only a few hundred nanometers. Besides that, the lateral electric field is strongest between the contacts. Thus, we are most likely probing a thin layer at the surface. From this simulation, a very sharp increase of the current would result; which actually is not the case in the presented experimental curves. The relaxation dynamics or cooling process is very fast and related to an extreme drop of the photocurrent within 1 \( \mu s \). Afterwards, it asymptotically approaches zero. Comparing the experimental results with the theoretical findings, it is clear that the amplitudes and time constants of the signals are different. Only in the case of the long relaxation process a thermal effect may contribute to the photocurrent signal. Furthermore, the simulation does
Figure 10. Simulation of a thermally induced photocurrent by a short laser pulse for five different temperatures as indicated. A steep current peak slowly relaxing to zero.

not reveal any current oscillations, which supports the view that instead the system is driven into a metastable condition different from a simple heating scenario.

For our case of (TMTTF)$_2$SbF$_6$, the laser photons with $\hbar\omega = 2.35$ eV excite either the transition from the quarter-filled HOMO band to the LUMO+1 band of the TMTTF$^+\delta$ molecules or the transition from a fully occupied band far below the Fermi energy to the quarter-filled unoccupied mid-infrared band, as depicted for the transitions labeled by I in Figure 14 of the Appendix A.2. After the vertical electronic excitation into the unoccupied bands, several decay channels are open:

- A free electron and hole can be created, which get separated from each other by the electric field, leading to the photocurrent signal. However, they can recombine before they reach the electrodes or become trapped. In general, the lifetime of free carries is between nanoseconds and microseconds.
- After a few picoseconds the excited electronic states couple to the lattice phonons and intramolecular vibrations. By this, energy is transferred to the lattice subsystem and the sample is warmed up.
- Afterwards, the excited electrons decay back into their initial state by emitting a red-shifted photon. The latter can be reabsorbed or detected as a photoluminescence signal, for instance in Raman scattering experiments. Also, a radiation-free decay is possible into the lowest excited (mid-infrared) band.
- Otherwise, the excited charge can relax into a low-lying unoccupied band, triggering a charge transfer from one cation to its neighboring molecule in the picture of a molecular chain. This can induce a phase transition from the charge-ordered to the metallic state.
- Recently, in photoluminescence spectra of a two-dimensional organic Mott insulator the creation of an exciton was suggested \[57\]. The theory of photo-induced transitions \[58\] proposes that higher excited states can be populated and excitons brake up; Frenkel excitons becomes charge-transfer excitons and propagate. The separation of excitons may also stimulate a light-induced transition. Especially in low-dimensional system fluctuations as well as electron-electron and electron-phonon interactions play a crucial role, eventually destabilizing the equilibrium state. This picture is similar to the model suggested for the photo-induced phase transition in TTF-CA \[56,59\].

While in Figs. 7 and 8 the laser pulse hits the sample after the voltage is applied, we now delay the voltage pulse with respect to the laser pulse; in Figure 11 the photocurrent signal is displayed for different delay times of the voltage pulse. For more than 0.8 ms delay time the signals overlap with each other. Interestingly, the oscillations persist even when the voltage pulse is applied long after the light irradiation; but with a smaller amplitude. Moreover, the oscillations appear still at the same position, which proves that they are rigidly linked to the light pulse. It is surprising that the current curves are strongly displaced with respect to each other, indicating that without an external electric field, the excited free electrons and holes (or excitons) are not separated from each other and
Figure 11. Raw data of the photocurrent signal measured for $(\text{TMTTF})_2\text{SbF}_6$ for various delay times between the voltage and 8 ns laser pulse for a laser energy of $E = 40 \mu J$ and $U = 10 \text{ V}$ at $T = 147 \text{ K}$.

hence, recombine rapidly. In the case of a pure thermally induced effect, the current dynamics or the resistance, respectively, is expected to be independent of the applied voltage and to exhibit the same temporal behavior, regardless of the delay time between laser pulse and voltage pulse.

At this point, we should recall similar effects detected over the last fifty years in other compounds with and without illumination.

1. In the 1960’s and 1970’s transport measurements on silicon with two Schottky contacts revealed current oscillations in the kHz regime, which were ascribed to double current injection at the contacts [60]. However, the frequency could be varied by changing the voltage bias or temperature and the oscillations occur only above a critical threshold field [61].

2. The well-known creation of hot or non-equilibrium electrons in semiconductors such as GaAs [32] leads to electric field-dependent charge carrier velocities, causing a negative-differential resistance. One of the related phenomena is the famous Gunn effect, yielding current oscillations in the microwave regime. In this case, the frequency can be tuned by the sample length or by the applied voltage.

3. Similar to the negative-differential resistance current oscillations in pure semiconductors, self-oscillating (photo-)currents were observed in superlattice structures, consisting of two different direct semiconductors such as GaAs-AlAs [62]. Due to the different energy levels in the quantum wells and the tunneling effect between the layers, a negative differential velocity regime is created as a function of the external electric field. The electric field reveals a spatial variation within the superlattice structure. Thereby, the different regions separated by domain walls get unstable and hence cause the oscillations. The resonance frequency can be tuned by the bias voltage or laser power [63,64].

4. Current oscillations occur frequently in nonlinear transport studies in various organic conductors [37,65–70]. There, they are often related to sliding charge-density waves (CDW) or electrically induced insulator-metal transitions accompanied by a bistability [46,65].

5. It is known that ferroelectric materials render a modification of the polarization by light stimulation, which can be detected by photoconductivity measurements [71]. Here, the detected signal is generally composed in of three different terms [72–74]: first, a fast decay of the excited states within a few microseconds, in addition a pyroelectric signal that can last several milliseconds, and finally piezoelectric oscillations superimposed on the decaying current due to the thickness variation of the sample by sound waves. Since the Fabre salts are charge-ordered
ferroelectric salts they resemble the behavior of well-known inorganic ferroelectric, such as LiNbO$_3$.

6. Above a certain threshold field, voltage oscillations (in the range of a few 100 mV/cm) occur in CDW systems [46,75], such as NbSe$_3$, TaS$_3$ or K$_{0.3}$MoO$_3$, which are attributed to the collective sliding of the charge density wave. There, the resonance frequency depends on the applied voltage and temperature. However, only very few photoconductivity studies were conducted on CDW materials [76,77]; most of them utilized a lock-in technique that does not directly provide the time-dependent behavior of the photocurrent. It was reported that the threshold voltage between the creeping and sliding state of the CDW can be raised by increasing the light intensity.

While the first four possibilities have already been disregarded as a possible origin of the current oscillations in our study, let us discuss the remaining suggestions. (TMTTF)$_2$SbF$_6$ can be considered as a one-dimensional charge-ordered ferroelectric, and the temporal dynamics of its photocurrent resembles that of illuminated inorganic ferroelectric crystals. For these, the current oscillations are related to the speed of sound in the bulk material. In our case a sample length of $l=2$ mm and a resonance frequency of 77.5 kHz would lead a sound velocity $v_{\text{sound}} = f \times l$ of 155 m/s. This is one order of magnitude below the speed of sound (2500 m/s$^{-1}$) previously reported [78–80]. Hence light-induced sound waves can be ruled out as origin of the observed oscillations.

The most appealing explanation of current oscillations in the charge-ordered state of (TMTTF)$_2$SbF$_6$ starts with a 4$k$-CDW. In general a coherent sliding of the density wave is observed above a more or less well defined electric threshold field of a few mV/cm [46,75]; this is not the case for (TMTTF)$_2$SbF$_6$. Here we have to apply an electric field of several V/cm, much more than typically needed for classical CDW systems; but two orders of magnitude smaller than the threshold fields in the charge-ordered salt $\alpha$-(BEDT-TTF)$_2$I$_3$ [22,28]. The large difference can be explained by the fact that most of the typical CDW compounds develop an incommensurate charge modulation that is pinned to impurities. Charge-ordered systems are inherently commensurate and much more rigid as far as the coupling to the underlying lattice is concerned. Furthermore, single-particle excitations should not weaken the pinning of the CDW, it may even increase the threshold field. In conventional CDW systems, photo excitation does not facilitate sliding of the condensate. In addition, it was observed in NbSe$_3$ and comparable CDW systems that the resonance frequency of the depinned 4$k$-CDW depends on the external voltage; in the present case of (TMTTF)$_2$SbF$_6$, however, no such effect could be identified. Further photoconductivity studies have to be performed on charge-ordered one-dimensional materials to clarify this point. Nevertheless, we suggest that our observations are caused by photo-induced fluctuations of the charge-ordered state, which alter the current flow of the single particles and hence, the photocurrent.

6. Summary and conclusion

We report on the first study of photoconductivity and non-linear transport in quasi-one-dimensional charge-ordered compounds. The Fabre salt (TMTTF)$_2$SbF$_6$ exhibits unusual conduction behavior when electric fields and light pulses are applied. Most interesting is the time-dependence of the photoconductivity signal which reveals current oscillations in the kHz range below the charge-order transition at $T_{\text{CO}} = 157$ K. Within milliseconds after photo excitation the signal exhibits an overall decay, in addition a faster decay process can be identified in the microsecond range. The most surprising observation, however, are oscillations superimposed on top of the photocurrent that contains several frequencies in the kHz regime. We investigate the field, laser power and temperature dependence. We discuss possible explanations of the observed photoconductivity and current oscillations; suggesting an interaction of the excited single particles and the charge-ordered state.
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Appendix A. Ab initio calculations

For a better understanding of the photoinduced processes in the Fabre salts, we have performed calculations of the electronic band structure and optical conductivity.

Appendix A.1. Band Structure

The band structure of $(\text{TMTTF})_2X$ (with $X=\text{PF}_6$, $\text{AsF}_6$, and $\text{SbF}_6$) were determined by ab initio Density Functional Theory (DFT) calculations which is implemented in the software package Quantum Espresso [81]. In contrast to the well-established Hückel theory with its molecular orbitals, it is based upon plane-waves and it is not semi-empirical; instead, it uses the advantage of the crystal periodicity and can be applied to almost every material. Moreover, it is fast, reliable and robust. Furthermore, we used a norm-conserving PBE GGA functional for all atom types, up to a certain level taking into account the exchange correlation as well as the spatial variation of the charge density. The room-temperature crystal structure for the calculations were taken for $(\text{TMTTF})_2\text{PF}_6$ and $(\text{TMTTF})_2\text{AsF}_6$, while for $(\text{TMTTF})_2\text{SbF}_6$ we performed x-ray scattering from $T = 300 \text{ K}$ down to $100 \text{ K}$ [13,14,82]; the structures were not relaxed within the calculations. Due to the stoichiometry of the organic salts, they are treated as metals. Thus, a smearing factor of 0.005 Ry was applied. Furthermore, the cutoff energy of the plane waves and the electron density were set to 30 Ry and 120 Ry, respectively. The wave functions were determined in a self-consistent way on a Monkhorst $8 \times 8 \times 4$ $k$ grid. The band structure was calculated on the $k$-path: $\Gamma (0,0,0) \rightarrow X (0.5,0,0) \rightarrow U (0.5,0.5,0) \rightarrow \Gamma (0,0,0) \rightarrow Y (0,0.5,0) \rightarrow T (0,0.5,0.5) \rightarrow \Gamma (0,0,0) \rightarrow Z (0,0.5) \rightarrow R (0.5,0.5,0.5) \rightarrow \Gamma (0,0,0)$.

The resulting band structure close to the Fermi energy is depicted in Figure 12 for $T = 300 \text{ K}$. The HOMO of the Fabre-salts is three quarter-filled. Due to the dimerization between the TMTTF$^+0.5$ molecules, the HOMO is split into a lower completely filled band and a half filled upper band. These bands are well separated from the other bands which are located at -2 eV below and about 1.5 eV above the HOMO band (not shown here). The anion layer does not contribute to the bands near the Fermi energy due to the closed shell of the anions. The occupied and unoccupied bands of the anions reside far below and above the HOMO band.

![Figure 12. Comparison of the band structure of (TMTTF)$_2X$, X= PF$_6$ (red), AsF$_6$ (blue), and SbF$_6$ (black) at room temperature along the k-path: $\Gamma (0,0,0) \rightarrow X (0.5,0,0) \rightarrow U (0.5,0.5,0) \rightarrow \Gamma (0,0,0) \rightarrow Y (0,0.5,0) \rightarrow T (0,0.5,0.5) \rightarrow \Gamma (0,0,0) \rightarrow Z (0,0.5) \rightarrow R (0.5,0.5,0.5) \rightarrow \Gamma (0,0,0)$.](image-url)
Figure 13. Band structure of \((\text{TMTTF})_2\text{SbF}_6\) in a temperature range from 300 K (dark red) to 100 K (dark blue) on the selected \(k\)-path: \(\Gamma(0,0,0) \rightarrow X(0.5,0,0) \rightarrow U(0.5,0.5,0) \rightarrow \Gamma(0,0,0) \rightarrow Y(0,0.5,0) \rightarrow T(0,0.5,0.5) \rightarrow \Gamma(0,0,0) \rightarrow Z(0,0.5,0) \rightarrow R(0.5,0.5,0.5) \rightarrow \Gamma(0,0,0)\).

The dimensionality of the compounds is reflected by the band dispersion. The band shape is similar for all compounds; the main deviations occur for the paths corresponding to the \(b\)-direction \((\Gamma \rightarrow Y\) or \(T \rightarrow \Gamma)\). They reveal an enhanced dispersion going from the large \(\text{SbF}_6^-\) to the smaller \(\text{PF}_6^-\) anions. This is consistent with the increase of the overlap between neighboring TMTTF molecules along the \(b\)-direction, being caused by the reduced spacing between the adjacent stacks due to the smaller anion size when going from \((\text{TMTTF})_2\text{SbF}_6\) to the smaller \((\text{TMTTF})_2\text{PF}_6\). Therefore, the compounds become more two-dimensional. The largest dispersion appears for the paths: \(\Gamma \rightarrow X\) which is along the stacking direction, \(U \rightarrow \Gamma\) and \(Z \rightarrow R\) containing contributions of the \(a\)-directions as well. The bandwidth for the three materials is about 350 meV for the upper as well as for the lower band, being in excellent agreement with previous studies \([13,14]\). For the \(b\)-direction, the bands are weakly warped and intersect the Fermi energy. In contrast, there is no energy dispersion for the \(c\)-direction \((\Gamma \rightarrow Z)\), since the adjacent cation chains are separated from each other along the \(c\)-axis by the anion layer. The dimerization gap at the \(X\)-point decreases from about 110 meV for \(\text{AsF}_6\) and \(\text{SbF}_6\) to 81 meV for \(\text{PF}_6\).

Moreover, also the temperature-dependence of the band structure was studied as visualized for \((\text{TMTTF})_2\text{SbF}_6\) over a temperature range between 300 K and 100 K in Figure 13. With decreasing temperature the bandwidth along the stacking direction \((\Gamma \rightarrow X)\) slightly decreases whereas it increases for the path \((\Gamma \rightarrow T)\) and \((T \rightarrow \Gamma)\). This implies that the compounds become more two-dimensional on cooling. This behavior is similar to applying chemical pressure, already discussed for Figure 12. This observation is supported by calculations of Jacko \textit{et al.} \([14]\) who also observed an increase of the dimensionality when the crystal structure of \((\text{TMTTF})_2\text{PF}_6\) under hydrostatic pressure was taken for the calculations. Concerning the dimerization, the corresponding gap at the \(X\)-point obviously diminishes from about 110 meV at 300 K to 55 meV at 100 K, again resembling the effect of pressure as demonstrated in Figure 12.

Appendix A.2. Optical spectra

Since optical spectra of the Fabre salts have not yet been theoretically investigated to our best knowledge, we have performed DFT calculations to derive the optical constants for all three crystallographic axes. While the excitation spectra were calculated for an evenly spaced grid of 132 \(k\)-points, the same functional, smearing factor, and cutoff energies were used as reported in the previous paragraph. The interband transitions were folded with a Lorentz function with a width of 100 meV, in contrast, for the Drude a width of 10 meV was assumed.
Figure 14. (a) Calculated optical conductivity of (TMTTF)$_2$SbF$_6$ between $T = 300$ and 100 K in a frequency range from 1000 cm$^{-1}$ to 30000 cm$^{-1}$ for the $a$, $b$, and $c$-direction. (b) Schematic diagram of the transitions between the occupied and unoccupied bands which are labeled by the corresponding markers from the optical conductivity spectra in (a).

Figure 14 displays the optical conductivity of (TMTTF)$_2$SbF$_6$ in a frequency range between 1000 cm$^{-1}$ and 30000 cm$^{-1}$ for temperatures between 300 K and 100 K. Below 5000 cm$^{-1}$, the $a$-direction is dominated by a Drude component (marked by I). While it is less pronounced in the $b$-direction, the $c$-direction does not reveal any Drude component and drops to zero as $\omega \rightarrow 0$. The ratio of the optical dc resistivity is $\rho_a : \rho_b : \rho_c = 1 : 10 : 6500$ which agrees well with the result of the transport measurements [10].

On top of the Drude feature several resonances II are superimposed which can also be found in the $c$-direction. These are excitations from the completely occupied band near the Fermi energy $E_F$ into the empty states above $E_F$, being the strongest for the $b$- and $c$-direction. Since no correlation effects are taken into account in the DFT calculations, the mid-infrared II band is not determined correctly. In the case of electron-electron correlations, the spectral weight of the Drude component is transferred to the mid-infrared band, leading to an enhancement of it. The weaker and broader interband transitions (III) occur for the $a$, $b$, and $c$-direction between 12 000 cm$^{-1}$ and 20 000 cm$^{-1}$. They can be attributed to the excitation from the half-occupied band to the lowest unoccupied band LUMO, which is located 1.7 eV above $E_F$, as depicted in Figure 14(b). Moreover, also the excitations from the fully occupied bands -2 eV below $E_F$ to the half-filled band contribute to this feature. The smeared feature IV resides above 20 000 cm$^{-1}$ and can be assigned to the transition from the HOMO to the unoccupied bands located at 2 eV in the band structure and from the occupied bands being located -2.5 eV below $E_F$ to the half-occupied band at $E_F$.

By comparing the intensity of the features III and IV for the various directions, it is evident that the bands in the stacking direction are by a factor of 10 less intense. This is because these transitions are caused by intramolecular electronic transitions of the cation molecule TMTTF$^{+0.5}$, whose dipole moment is along the molecular axis. From the crystal structure it can be deduced that the long molecular axis of the TMTTF molecules is mainly oriented along the $c$-axis. For instance, the transition from the occupied band at the Fermi energy to the band at 1.7 eV corresponds to the
transition from the HOMO of the TMTTF$^{+0.5}$ molecule in the LUMO. The excitation energy of this transition decreases with increasing molecule size.

Besides the overall shape of the spectra and the assignment of the features it is very interesting to track the temperature evolution of the spectra and compare it to the experimental measurements. The II band reveals the largest temperature dependence. In the $a$-direction, its intensity decreases on cooling and vanishes in the Drude component, while for the $b$- and $c$-direction it only shifts to lower and to higher frequencies, respectively. In addition, also the III- and IV-bands exhibit a temperature-dependent modification. Therefore, also the electronic excitations into higher energy levels can be used to track thermal effects. This is also true for phase transitions which are accompanied by a strong transfer of spectral weight.

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