Research Article

Double-Grating Displacement Structure for Improving the Light Extraction Efficiency of LEDs

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To improve the light extraction efficiency of light-emitting diodes (LEDs), grating patterns were etched on GaN and silver film surfaces. The grating-patterned surface etching enabled the establishment of an LED model with a double-grating displacement structure that is based on the surface plasmon resonance principle. A numerical simulation was conducted using the finite difference time domain method. The influence of different grating periods for GaN surface and silver film thickness on light extraction efficiency was analyzed. The light extraction efficiency of LEDs was highest when the grating period satisfied grating coupling conditions. The wavelength of the highest value was also close to the light wavelength of the medium. The plasmon resonance frequencies on both sides of the silver film were affected by silver film thickness. With increasing film thickness, plasmon resonance frequency tended toward the same value and light extraction efficiency reached its maximum. When the grating period for the GaN surface was 365 nm and the silver film thickness was 390 nm, light extraction efficiency reached a maximum of 55%.

1. Introduction

Low-carbon economy and green economy have recently become major themes of global development. As a type of green lighting energy, light-emitting diodes (LEDs) have attracted considerable research attention [1]. LEDs present the advantages of energy conservation, environment friendliness, fast response, and long lifetimes. Thus, LEDs lighting systems have been considered to reduce greenhouse gas emissions by replacing fuel-based lighting in the developing world [2]. Not only that, these light sources are extensively used as backlight sources of liquid crystal displays, advertisements, night lights, traffic signals, outdoor lighting, full color displays [3], and biometric devices [4], in addition, the array of LEDs has been used to inactivate bacteria in liquid suspension and on exposed surfaces [5], among other applications. Despite these advantages, the development of LED core technology is constrained by the photons produced in the lighting devices. Such a constraint prevents photons from being effectively radiated and converted into available optical power. This problem stems from the inherent structure of solid-state lighting devices: the refractive index of solid semiconductor luminophore is higher than that of the surrounding medium [6]. Thus, the light emitted by the lighting devices easily causes total internal reflection in the interface of semiconductor and air, after which the light retraces the optical path of the luminophore and is converted into heat [7]. This phenomenon makes the LEDs working in high-temperature environments for a long time; it not only wastes energy, but also shortens LED lifetimes. So the improvement of LEDs light extraction efficiency can not only extend lifetime, but also make important contribution to the development of the new green energy.

Foreign and domestic experts have begun improving the structure of LED chips, thereby significantly enhancing light extraction efficiency. These improvements primarily include changing the geometric shape of the chips, reducing the duration of total internal reflection in the chips, and decreasing light energy losses. The methods generally applied in realizing such enhancements are using the inverted pyramid structure [8, 9], performing surface roughening [10], constructing a mirror with Bragg grating reflection [11], and creating a two-dimensional photonic crystal structure on the surface of GaN [12–15]. Using the characteristic of surface
2. Theories and Methods

2.1. Introduction to SP Resonance. When the wave frequency for the excitation of SPs is larger than the plasmon resonance frequency ($\omega_p$), the electromagnetic field is radiatively propagated far from the space between the metal/dielectric material interface, a phenomenon described as the radiative mode of SPs [20, 21]. Only when the horizontal direction wave vector $k_x$ is very small (i.e., resonance frequency $\approx \omega_p$) can the life cycle of the radiant mode of SPs on a metal surface be defined. The gradual increase in $k_x$ results in a resonance mode life cycle that is too short to have any practical significance. The SP resonance in the range of visible frequency is the nonradiative SP mode. The electromagnetic field produced by this mode is limited to an area near the metal surface. This phenomenon is called electromagnetic wave dispersion.

The SP dispersion relationship is depicted in Figure 1. The dispersion curve of the nonradiative SPs produced on a metal surface lies completely on the right side of the incident electromagnetic wave line of the dielectric material [20]. This result indicates that the SPs have a longer horizontal component of wave vector $k_x$ than that of the incident electromagnetic wave vector $k_0$ of the same energy. Thus, the general electromagnetic waves incident from the dielectric material cannot stimulate nonradiative SPs and reach the resonance mode. To match the excitation conditions for the SP resonance mode, a coupled mechanism should be applied. This mechanism increases wave vector $\Delta k_x$ and causes the incident electromagnetic wave to achieve a higher $k_x$ (Figure 2).

Two coupled mechanisms are commonly used [20]. One involves producing a small-raster periodic structure on
a metal surface as a coupled medium. Under an external electromagnetic field, the free electrons of the metal surface in these cycle structures cause polarized electron oscillation for a specific wavelength. The generated electromagnetic field can provide an additional value $\Delta k$ to the incident electromagnetic wave. Subsequently, SP resonance is stimulated. The other coupled mechanism uses the total internal reflection dissipation field produced in a material with a high dielectric constant as the excitation source to stimulate SP resonance.

The electromagnetic wave vector can be written as

$$k^2 = \frac{\omega^2 \varepsilon \mu}{c^2}, \quad (1)$$

where $\varepsilon$ and $\mu$ are the relative permittivity and relative permeability of the material, respectively. When the electromagnetic wave passes through the prism of high permittivity, the wave vector is greater than that observed in dielectric material $\varepsilon_1$. When the incident light produces total internal reflection in the prism/dielectric material interface, part of the evanescent field is tunneled into the dielectric material near the total internal reflection interface. The wave vector $k$, of the evanescent field and the total internal reflection wave vector have the same size when the distance between the prism and the metal surface is sufficiently small. At the same time, when the incident light wave vector horizontal component of total internal reflection satisfies the conditions of SPs, the SPs in the dielectric material/metal interface are stimulated.

2.2. Finite Difference Time Domain Method. The finite difference time domain (FDTD) method [22], an extension of the finite difference method, is a numerical analysis approach that directly performs computer simulations by using Maxwell curl equations for electromagnetic fields. Using the Maxwell curl equations in accordance with the field quantities that have special configurations in the Yee grid yields [23, 24]

$$\frac{\varepsilon (i + 1/2, j, k)}{\Delta t} [E_x^{p+1}(i + 1/2, j, k) - E_x^p(i + 1/2, j, k)]$$

$$+ \sigma(i + 1/2, j, k)[E_x^{p+1}(i + 1/2, j, k) + E_x^p(i + 1/2, j, k)]$$

$$= \frac{[H_y^{p+1/2}(i+1/2, j+1/2, k) - H_y^{p+1/2}(i+1/2, j-1/2, k)]}{\Delta y}$$

$$- \frac{[H_y^{p+1/2}(i+1/2, j, k+1/2) - H_y^{p+1/2}(i+1/2, j, k-1/2)]}{\Delta z}, \quad (2)$$

Let

$$p = \left( i + \frac{1}{2}, j, k \right), \quad q^+ = \left( i + \frac{1}{2}, j + \frac{1}{2}, k \right),$$

$$q^- = \left( i + \frac{1}{2}, j - \frac{1}{2}, k \right), \quad (3)$$

$$r^+ = \left( i + \frac{1}{2}, j, k + \frac{1}{2} \right), \quad r^- = \left( i + \frac{1}{2}, j, k - \frac{1}{2} \right).$$

Thus,

$$E_x^{p+1} = \left[ 1 - \Delta t \sigma p / 2 \varepsilon_p \right] E_x^p + \left( \frac{\Delta t / \varepsilon_p}{1 + \Delta t \sigma p / 2 \varepsilon_p} \right) \left[ \frac{[H_y^{p+1/2} - (H_y^{p+1/2})] / \Delta y}{\Delta y} - \frac{[H_y^{p+1/2} - (H_y^{p+1/2})] / \Delta z}{\Delta z} \right], \quad (4)$$

where the superscripts $n + 1/2$ and $n$ are the time step-numbers; a group of changes $(i, j, k)$ and $\Delta t$ are the time steps; $\Delta x, \Delta y, \Delta z$ are the distances of the adjacent lattice point in the $x, y, z$ directions, respectively. Other field quantities can be treated in the same manner.

To perform a stable numerical simulation, the variables should satisfy

$$v_{\text{max}} \Delta t \leq \frac{1}{\sqrt{(1/\Delta x)^2 + (1/\Delta y)^2 + (1/\Delta z)^2}}, \quad (5)$$

where $v_{\text{max}}$ is the maximum phase velocity, $\Delta t$ denotes the time step, and $\Delta x, \Delta y, \Delta z$ are the distance steps in the $x, y, z$ directions, respectively. When $\Delta x = \Delta y = \Delta z = \Delta$, (5) can be simplified as $v_{\text{max}} \Delta t < \Delta / \sqrt{3}$. To reduce the error caused by numerical dispersion effects, (5) should also satisfy the conditions $\Delta / \lambda_{\text{min}} < 1/10$, where $\lambda_{\text{min}}$ is considered the shortest wavelength of the electromagnetic wave considered.

The FDTD method was used to solve the numerical analysis problems presented by grating-patterned etching on the metallic silver film. This approach also enables the excitation of SPs and application of the total internal reflection dissipation field for exciting the SPs. The changes in time step $\Delta t$ reflect the variations in energy flow on the receiving surface. This reflection is realized by placing a receiving surface over the structure. The changes in energy flow facilitate the analysis of the effects of structural parameter variations on the light extraction efficiency of LEDs.

3. Physical LED Model

The physical model of the LED with a double-grating displacement structure was constructed by grating-patterned etching on the GaN and metallic silver film surfaces (Figure 3). The model comprises the silver film layer, P-GaN layer, active layer, N-GaN layer, and SiC substrate. During the experiment, the calculated area dimensions of 4200 nm × 420 nm were realized; the thicknesses of the P-GaN and N-GaN layers were 200 and 400 nm, respectively. Compare with the physical LED model as described in [17], a grating-patterned etching was added to the GaN surface, which had two reasons. (1) Grating-patterned etching can diminish part of the total internal reflection produced by incident light. (2) A portion of the light causes the evanescent field of total internal reflection to excite SPs; The study in [7] has discussed that grating-patterned etching on the top surface can effectively enhance the light extraction efficiency of LEDs; thus, this phenomenon results in a combined effect with
4. Numerical Simulation and Analysis

A numerical simulation was conducted using the FDTD method. The wavelength of light was 500 nm, the refractive index of air was 1.0, and the refractive index of GaN was 2.4. The silver film was used in the Lorentz dispersion model. In analyzing the light absorption process, the silver film is disregarded. The light emitted by the active layer is replaced by the total field-scattered field source, an approach that realizes incident plane waves from all directions. The intensity was 1 and the grating period of the metallic silver film was 375 nm. A receiving surface, which is capable of reflecting light intensity, was placed above the model to keep track of energy flow.

4.1. Effect of Grating Period on the Light Extraction Efficiency of the GaN Surface. Figure 4 presents the energy flow of the GaN with a double-grating displacement structure at different grating periods. The maximum flow intensity value was 0.44, which can be obtained at a GaN surface grating period of 365 nm. Figure 5 shows the relationship between the surface grating period of GaN and light extraction efficiency when the silver film thickness changes from 300 nm to 400 nm. For light extraction efficiency with oscillatory changes in grating period, the maximum value acquired constantly ranged from 360 nm to 380 nm. This result confirms that grating period influences the light extraction efficiency of LEDs. At a silver film thickness of 300 nm to 400 nm, the relationship curves of grating period and light extraction efficiency are similar and the maximum light extraction efficiency occurred at a 365 nm grating period. When the grating period continued to increase, the light extraction efficiency gradually declined. This result is attributed to two factors. First, the surface grating period is related to the wave vector that causes plasmon resonance. Thus, only when the grating period satisfies the excitation conditions can the SPs be excited. Consequently, visible light is extracted and transmissible light is enhanced. Second, the ideal size of diffraction grating satisfies the wavelength of light on a metal surface when the period is considerably smaller than the wavelength of light. Otherwise, conductivity to light diffraction is low and the effect of changing light directions caused by light diffraction decreases. On the other hand, when the period is far from the wavelength of light, the plane of the grating section is excessively large, thereby causing stronger total internal reflection, which slows down light emission.

4.2. Effect of Silver Film Thickness on Light Extraction Efficiency. The grating period was set at 365 nm on the basis of the analysis in Section 4.1. The thickness of the metallic silver film was set at 300 nm to 400 nm. Figures 6 and 7 show the energy flow of the silver film with a double-grating displacement structure and that of the silver film with a single grating, respectively. The light extraction efficiency of LEDs with a double-grating displacement structure improved by 43% over that of the silver film surface with a single grating. The graph shows that the silver film thickness should satisfy the grating coupling conditions to enhance light extraction efficiency. A silver film that is too thick or too thin cannot induce coupling. The plasmon resonance mode existed on both sides of the metal/dielectric material interface. Thus, when the film thickness was at the nanometer scale, the evanescent field formed by the SPs was strong enough to pass through the other side of the metal. Consequently, the SP electromagnetic fields on both sides of the metal film interacted with each other and formed a group of coupled SPs, as illustrated by the energy flow curve of silver thickness at 390 nm in Figure 6.

The plasmon resonance frequency reached the same value and energy flow intensity reached a maximum of 0.55.
5. Conclusion

Grating-patterned etching was performed on GaN and silver film surfaces to establish an LED model with a double-grating displacement structure. The FDTD method was used in conducting a numerical simulation. A grating etched on the GaN surface effectively extracts a limited amount of light and enhances the light extraction efficiency of LEDs. At a GaN grating period of 365 nm and a silver film thickness of 390 nm, light extraction efficiency reaches a maximum of 55%.

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References

[1] X. Gu, T. Qiu, W. Zhang, and P. K. Chu, “Light-emitting diodes enhanced by localized surface plasmon resonance,” Nanoscale Research Letters, vol. 6, no. 1, article 199, 2011.

[2] E. Mills and A. Jacobson, “From carbon to light: a new framework for estimating greenhouse gas emissions reductions from replacing fuel-based lighting with LED systems,” Energy Efficiency, vol. 4, no. 4, pp. 523–546, 2011.

[3] H. Lin, S. Liu, X. S. Zhang, B. L. Liu, and X. C. Ren, “Enhanced external quantum efficiency of light emitting diodes by fabricating two-dimensional photonic crystal sapphire substrate with holographic technique,” Acta Physica Sinica, vol. 58, no. 2, pp. 959–963, 2009.
[4] N. Kourkoumelis and M. Tzaphlidou, “Eye safety related to near infrared radiation exposure to biometric devices,” The Scientific World Journal, vol. 11, pp. 520–528, 2011.

[5] L. E. Murdoch, M. MacLean, E. Endarko, S. J. MacGregor, and J. G. Anderson, “Bactericidal effects of 405nm light exposure demonstrated by inactivation of escherichia, salmonella, Shigella, Listeria, and mycobacterium species in liquid suspensions and on exposed surfaces,” The Scientific World Journal, vol. 2012, Article ID 137805, 8 pages, 2012.

[6] D. H. Kim, C. O. Cho, Y. G. Roh et al., “Enhanced light extraction from GaN-based light-emitting diodes with holographically generated two-dimensional photonic crystal patterns,” Applied Physics Letters, vol. 87, no. 20, Article ID 203508, 3 pages, 2005.

[7] S. Trieu, X. Jin, A. Ellaboudy et al., “Top transmission grating GaN LED simulations for light extraction improvement,” in Physics and Simulation of Optoelectronic Devices XIX, vol. 7933 of Proceedings of SPIE, San Francisco, Calif, USA, January 2011.

[8] C. M. Tsai, J. K. Sheu, W. C. Lai et al., “Enhanced output power in GaN-based LEDs with naturally textured surface grown by MOCVD,” IEEE Electron Device Letters, vol. 26, no. 7, pp. 464–466, 2005.

[9] W. Zhou, G. Min, Z. Song, J. Zhang, Y. Liu, and J. Zhang, “Enhanced efficiency of light emitting diodes with a nano-patterned gallium nitride surface realized by soft UV nanoimprint lithography,” Nanotechnology, vol. 21, no. 20, Article ID 205304, 2010.

[10] K. S. Kim, S. M. Kim, H. Jeong, M. S. Jeong, and G. Y. Jung, “Enhancement of light extraction through the wave-guiding effect of ZnO sub-microrods in InGaN blue light-emitting diodes,” Advanced Functional Materials, vol. 20, no. 7, pp. 1076–1082, 2010.

[11] C. M. Lim and G. Hugh Song, “Design of superperiodic photonic-crystal light-emitting plates with highly directive luminance characteristics,” Journal of the Optical Society of America B, vol. 26, no. 2, pp. 328–336, 2009.

[12] A. David, H. Benisty, and C. Weisbuch, “Optimization of light-diffracting photonic-crystals for high extraction efficiency LEDs,” IEEE/OSA Journal of Display Technology, vol. 3, no. 2, pp. 133–148, 2007.

[13] D. H. Long, I. K. Hwang, and S. W. Ryu, “Design optimization of photonic crystal structure for improved light extraction of GaN LED,” IEEE Journal on Selected Topics in Quantum Electronics, vol. 15, no. 4, pp. 1257–1263, 2009.

[14] A. I. Zhmakin, “Enhancement of light extraction from light emitting diodes,” Physics Reports, vol. 498, no. 4-5, pp. 189–241, 2011.

[15] T. A. Truong, L. M. Campos, E. Matioli et al., “Light extraction from GaN-based light emitting diode structures with a noninvasive two-dimensional photonic crystal,” Applied Physics Letters, vol. 94, no. 2, Article ID 023101, 33 pages, 2009.

[16] V. K. Komarala, W. H. Guo, and M. Xiao, “Surface plasmon density of states at the metal-dielectric interface: dependence of metal layer thickness and dielectric material,” Journal of Applied Physics, vol. 107, no. 1, Article ID 014309, 5 pages, 2010.

[17] Y. Z. Lin, K. Li, F. M. Kong, J. Zhao, L. G. Du, and H. Gao, “Comprehensive numeric study of gallium nitride light-emitting diodes adopting surface-plasmon-mediated light emission technique,” IEEE Journal on Selected Topics in Quantum Electronics, vol. 17, no. 4, pp. 942–951, 2010.

[18] Z. Wu, J. W. Haus, Q. Zhan, and R. L. Nelson, “Plasmonic notch filter design based on long-range surface plasmon excitation along metal grating,” Plasronics, vol. 3, no. 2-3, pp. 103–108, 2008.

[19] S. H. Choi, S. J. Kim, and K. M. Byun, “Characteristics of light emission from surface plasmons based on rectangular silver gratings,” Optics Communications, vol. 283, no. 14, pp. 2961–2966, 2010.

[20] H. Raether, Surface Plasmons on Smooth and Rough Surface and on Gratings, Springer, Berlin, Germany, 1998.

[21] S. Glasberg, A. Sharon, D. Rosenblatt, and A. A. Friesem, “Long-range surface plasmon resonances in grating-waveguide structures,” Applied Physics Letters, vol. 70, no. 10, pp. 1210–1212, 1997.

[22] K. S. Yee and J. S. Chen, “The finite-difference time-domain (FDTD) and the finite-volume time-domain (FVTD) methods in solving Maxwell’s equations,” IEEE Transactions on Antennas and Propagation, vol. 45, no. 3, pp. 354–363, 1997.

[23] X. Zhou and W. Lin, “Numerical stability and numerical dispersion of compact conformal mapping 2D-FDTD method used for the dispersion analysis of waveguides,” International Journal of Infrared and Millimeter Waves, vol. 20, no. 7, pp. 1403–1412, 1999.

[24] M. T. Bettencourt, “Flux limiting embedded boundary technique for electromagnetic FDTD,” Journal of Computational Physics, vol. 227, no. 6, pp. 3141–3158, 2008.