Experimental Validation of Low-Z Ion-Stopping Formalisms around the Bragg Peak in High-Energy-Density Plasmas

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We report on the first accurate validation of low-Z ion-stopping formalisms in the regime ranging from low-velocity ion stopping—through the Bragg peak—to high-velocity ion stopping in well-characterized high-energy-density plasmas. These measurements were executed at electron temperatures and number densities in the range of 1.4–2.8 keV and 4 × 1023–8 × 1023 cm−3, respectively. For these conditions, it is experimentally demonstrated that the Brown-Preston-Singleton formalism provides a better description of the ion stopping than other formalisms around the Bragg peak, except for the ion stopping at vi ∼ 0.3vth, where the Brown-Preston-Singleton formalism significantly underpredicts the observation. It is postulated that the inclusion of nuclear-elastic scattering, and possibly coupled modes of the plasma ions, in the modeling of the ion-ion interaction may explain the discrepancy of ∼20% at this velocity, which would have an impact on our understanding of the alpha energy deposition and heating of the fuel ions, and thus reduce the ignition threshold in an ignition experiment.

In hot-spot ignition experiments [1] at the National Ignition Facility [2], which use deuterium-tritium (DT) fuel, an understanding of the DT-alpha energy deposition and heating of the high-energy-density (HED) plasma is critical for determining the ignition threshold. This requires a fundamental understanding of the DT-alpha stopping around the Bragg peak, where the ion velocity (vi) is similar to the average velocity (vth) of the thermal plasma electrons, for a wide range of electron (Te) and ion temperatures (Ti) and electron-number densities (ne) [3]. Ion stopping in HED plasmas has therefore been subject to extensive analytical and numerical studies for decades [4–14], but a theoretical treatment of ion stopping, especially around the Bragg peak, remains a difficult problem. The consensus is that the ion stopping at vi ∼ vth is treated well by the Born approximation [12] because the interaction between the fast ions and the plasma electrons is small, resulting in small energy transfers compared to the kinetic energy of the ions. At vi < vth, the ion stopping is harder to characterize but generally described by collisional theories that treat two-body collisions and large-angle scattering between the ions and the plasma electrons [13,15]. At ion velocities near vth, the Born approximation breaks down because scattering is no longer weak and collisional theories have difficulty providing a complete, self-consistent picture of the ion stopping due to the dynamic dielectric response of the plasma electrons. Rigorous quantum mechanical treatments based on convergent kinetic theories [6,14,16] try to rectify these challenges by utilizing the strengths of the different approaches applied to the different regimes; however, it is not clear how best to combine them and quantify their errors. Precise measurements of the ion stopping around the Bragg peak are therefore essential to guiding the theoretical efforts.

Although numerous efforts have been made to theoretically describe the behavior of ion stopping in HED plasmas, only a limited set of experimental data exists to test these theories. In addition, most of these experiments used only one particle with a distinct velocity in the high-velocity ion-stopping regime (vi > vth) [17–30] and thus did not simultaneously probe the detailed characteristics of the Bragg peak below and above vth. To the best of our knowledge, only two experiments have made an attempt to simultaneously probe the low- and high-velocity sides of the Bragg peak. The first experiment was conducted by Hicks et al. [28], who provided a qualitative description of the ion stopping around the Bragg peak. The second one was conducted by Frenje et al. [29], who provided the first experimental evidence that the position and magnitude of
TABLE I. Experimental parameters and key HED-plasma parameters.

| Shot  | Capsule* | Laser Peak Power (kJ) | Laser Gas Pressure (atm) | DHe Bang Times [ps]** | DD Bang Times [ps]** | DHe Yield | DD Yield | X-ray Yield | Ti (keV) | Te (keV) | SiO2 (keV) |
|-------|----------|-----------------------|--------------------------|-----------------------|---------------------|-----------|----------|-------------|----------|----------|------------|
| 75694 | 3He:6.81D3:(3)Ar:0.15S02:27 | 12.0 | 1.5 × 10^10 | 114.0 | 88 ± 0.3 | 1.5 × 10^10 | 114.0 | 88 ± 0.3 | 1.5 × 10^10 | 114.0 | 88 ± 0.3 | 1.5 × 10^10 |
| 75703 | 3He:6.71D3:(3)Ar:0.14S02:27 | 11.7 | 1.4 × 10^10 | 118.0 | 83 ± 0.3 | 1.4 × 10^10 | 118.0 | 83 ± 0.3 | 1.4 × 10^10 | 118.0 | 83 ± 0.3 | 1.4 × 10^10 |
| 75705 | 3He:6.61D3:(3)Ar:0.13S02:27 | 9.9 | 1.1 × 10^10 | 121.5 | 78 ± 0.3 | 1.1 × 10^10 | 121.5 | 78 ± 0.3 | 1.1 × 10^10 | 121.5 | 78 ± 0.3 | 1.1 × 10^10 |
| 75697 | 3He:6.61D3:(3)Ar:0.13S02:27 | 9.9 | 1.1 × 10^10 | 119.5 | 76 ± 0.3 | 1.1 × 10^10 | 119.5 | 76 ± 0.3 | 1.1 × 10^10 | 119.5 | 76 ± 0.3 | 1.1 × 10^10 |
| 75699 | 3He:6.61D3:(3)Ar:0.13S02:27 | 8.1 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 |
| 75700 | 3He:6.61D3:(3)Ar:0.13S02:27 | 8.1 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 |
| 75701 | 3He:6.61D3:(3)Ar:0.13S02:27 | 8.1 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 |
| 75702 | 3He:6.61D3:(3)Ar:0.13S02:27 | 8.1 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 | 120.5 | 66 ± 0.3 | 9.9 × 10^9 |

*Capsule thicknesses (in μm). The capsule diameter is 1 mm. The capsule X-ray temporal diagnostic.[33]

**X-ray times (at time 20–80 ns using the Particle X-ray Temporal Diagnostic.[33])

The experiments reported herein were carried out at OMEGA,[31] where eight deuterium-helium-3 gas-filled capsules were symmetrically imploded with 60 laser beams, delivering up to 12.0 kJ to the capsule in a 1-ns square pulse. As shown in Table I, each SiO2 capsule had a shell thickness of ~2.7 μm and an initial DHe-gas-fill pressure in the range of 12 to 13 atm. These capsules were also filled with a small amount of argon for a time- and space-resolved measurement of the electron-temperature and electron-number-density profiles.[32]

Essential to this Letter is to accurately characterize the spatially and temporally varying HED-plasma conditions during the nuclear-production period. For this, an unprecedented set of complementary nuclear and x-ray measurements was conducted, as illustrated in Table I and Figs. 1 and 2. Table I shows measured nuclear and x-ray bang times, burn-averaged \( T_\alpha \), and DD and DHe yields, and \( T_\alpha \) at the center of the implosion for all shots. Implosion parameters inferred from the measured data, essential to the ion-stopping predictions, are also shown in Table I. It is also notable that each shot pair at a given laser energy is producing reproducible data. Figure 1 shows the measured and modeled x-ray-emission history, and the DD-burn and DHe-burn histories together with the implosion trajectory for shot 75699, while Fig. 2 shows measured electron-number-density and electron-temperature profiles.

FIG. 1. Measured (points) and modeled (solid curves) nuclear-burn and x-ray-emission histories, and implosion trajectory for shot 75699. The x-ray-emission history was determined from x rays measured in the energy range of 3.375–3.600 keV. The implosion trajectory, which is well modeled by a 1D HYADES simulation, was measured with a time-gated imaging camera probing soft x rays from the SiO2 shell.
From Table I and Fig. 2, it is clear that the 1D-simulated $T_e$ profiles at the center of the implosion agree well with the measured $T_{e,0}$ value, which raises our confidence that the measured and inferred implosion parameters used for the ion-stopping predictions are determined with high accuracy (~10% considering all measurements and modeling). As illustrated in Table I, it is also notable that the burn-averaged $T_i$ values are significantly higher than the measured $T_e$. The reason for this is that the converging shock predominantly transfers energy to the heavier ions in the HED plasma. As the shock rebounds at the center of the implosion, it significantly raises $T_i$ and initiates the DD and D$^3$He nuclear reactions. Given that the ion-ion equilibration time is ~50 ps for these HED-plasma conditions, the ions are not fully in thermal equilibrium at the end of the ~170 ps long burn, and as a consequence the neutron-time-of-flight (NTOF)-measured values in Table I represent an apparent $T_i$. In addition, as the ions and electrons do not have time to fully equilibrate during burn (the ion-electron thermalization time is ~500 ps), the measured $T_e$ is consequently lower than the measured apparent $T_i$. By contrast, the electron-electron thermalization time is subpicosecond for these conditions, which implies that the electrons are internally in thermal equilibrium and are well described by the HYADES simulations. From a burn-averaged point of view, assigning $T_e$ to these plasmas is therefore meaningful.

For accurate experimental validation of the ion stopping around the Bragg peak, the energy losses ($-\Delta E_i$) of DD tritons (DD-$t$), DD protons (DD-$p$), D$^3$He alphas (D$^3$He-$\alpha$) and D$^3$He protons (D$^3$He-$p$), while traversing the well-characterized HED-plasma conditions, were simultaneously measured. An example of measured spectra of DD-$t$, DD-$p$, D$^3$He-$\alpha$, and D$^3$He-$p$ is shown in Fig. 3 for shot 75699 (see the detailed discussion about these measurements and the associated uncertainties in the Supplemental Material [37]). These spectra were measured with a single spectrometer, but other spectrometers fielded around the implosion were also used for these measurements [38]. The vertical arrows in Fig. 3 indicate the median energy for each measured spectrum, and by contrasting these energies to the average-birth energies (vertical dashed lines), $-\Delta E_i$ was determined to an accuracy of ~10% (see the Supplemental Material [37]) and used for the assessment of the ion stopping in the HED plasma. As the fusion products traverse the HED plasma with varying electron temperatures and electron-number densities (see Fig. 2), they probe the ion stopping predictions are determined with high accuracy (considering all measurements and modeling). As $-\Delta E_i$ values are significantly higher than the measured $T_e$ values for the low-velocity fusion products (DD-$t$, D$^3$He-$\alpha$, and DD-$p$), $-\Delta E_i$ values for all shots, where the measured $-\Delta E_i$ for the low-velocity fusion products were corrected for based on the different $dE/dx$-weighted $\langle T_e \rangle$, where the $-\Delta E_i$ value for the D$^3$He-$p$ was corrected for the burn-averaged $\langle n_eL \rangle$ change from the

FIG. 2. Profiles of (a) electron-temperature and (b) electron-number density for shot 75699 electron-temperature and electron-number density for shot 75699, measured by the MMI (data points) and simulated by HYADES (solid curves). These profiles were integrated over a time window of 1.23–1.33 ns and are x-ray emissivity weighted towards the end of the time window. To match the measured profiles, scaling factors in the ranges of 1.2–1.6 and 0.7–0.9 were applied to the HYADES-simulated electron-temperature and electron-number-density profiles for the eight shots, respectively.

FIG. 3. Measured spectra of DD-$t$, D$^3$He-$\alpha$, DD-$p$, and D$^3$He-$p$ for shot 75699. These fusion products are produced by the reactions $D+D \rightarrow (1.01\text{MeV}) + p (3.02\text{MeV})$ and $D+^3\text{He} \rightarrow ^4\text{He} (3.71\text{MeV}) + p (14.63\text{MeV})$, where the energies in the parentheses are the fusion-product birth energies (at zero ion temperature).
D$_3$He-BT to DD-BT, as discussed in the Supplemental Material [37]. The solid curves in Fig. 4 were obtained by integrating the Brown-Preston-Singleton (BPS) plasma-stopping-power function, describing only the ion-electron Coulomb interaction, for the $dE_i/dx$-weighted $\langle T_e \rangle$ values shown in Table I. Clearly, the data demonstrate that the BPS formalism is providing a good description of the ion stopping for these HED-plasma conditions, except for the stopping of DD-t at $v_i \sim 0.3v_{th}$. At this velocity, the BPS formalism systematically underpredicts DD-t energy loss for all shots. A systematic error in the measured DD-t energy loss can be ruled out in explaining this observation, as similar systematic errors would be evident in the measured D$_3$He-α and DD-p energy loss. An ion-bulk flow of $\sim 500$ km/s systematically in the direction away from the spectrometer, necessary to explain the observation, can also be excluded because spectrometers with nearly orthogonal lines of sight observe similar energy loss, and it would also be evident in the measured D$_3$He-α and DD-p spectra. A systematically too high DD-t $dE_i/dx$-weighted $\langle T_e \rangle$ for all shots can also be ruled out because a 300–400 eV lower value is required to explain the data, which is not plausible.

To examine the different ion-stopping formalisms routinely used in the field of inertial confinement fusion and to illustrate different approaches in unifying the different physical processes that dictate the characteristics of the Bragg peak, Fig. 5 contrasts the energy-loss data with predictions by BPS and Li and Petrasso (LP) [7] for shot 75699. As shown by the comparison, the BPS formalism is providing a better description of the Bragg peak, which supports the general view that the BPS formalism is considered to more accurately unify the binary-collision and dielectric-response formalisms with more rigorous quantum-diffraction corrections to the total ion stopping at $v_i \sim v_{th}$. On the other hand, the BPS formalism, considering only ion-electron Coulomb interactions, systematically underpredicts the DD-t energy loss at $v_i \sim 0.3v_{th}$ for all shots. This observation cannot be explained by the inclusion of ion-ion Coulomb scattering in the modeling because ion-stopping theories based on multi-ion-component responses predict that the contribution of the ion-ion Coulomb scattering to the total DD-t plasma-stopping power is $\sim 10\%$ at $v_i \sim 0.3v_{th}$ [40] (see the dashed curve in the inset of Fig. 5). This points to the idea that the contribution from the ion-ion component to the total ion stopping at this velocity could in fact be larger than predicted by the theories. This is certainly plausible, as all theories ignore the ion-ion nuclear-elastic scattering, which is more strongly weighted towards large-angle scattering than Coulomb scattering. To explain the data at $v_i \sim 0.3v_{th}$, the total ion stopping must be increased by $\sim 20\%$ (see the dotted curve in the inset of Fig. 5), possibly due to ion-ion nuclear-elastic scattering [41]. This postulation, if correct, would have an impact on our understanding of DT-alpha heating of the fuel ions in an ignition experiment. Another possibility that must also be considered in explaining this discrepancy is that coupled modes of the plasma ions are not considered in these theories. However, this is unlikely, as the ion-ion coupling is weak.

In summary, ion stopping around the Bragg peak has been measured in well-characterized HED-plasma conditions. This effort significantly advances previous efforts by providing the first accurate experimental validation of ion-stopping formalisms in the regime ranging from low-velocity ion stopping—through the Bragg peak—to high-velocity ion stopping. The data indicate that the BPS formalism provides a better description of the ion stopping than other formalisms around the Bragg peak, except for the ion stopping at
$v_i \sim 0.3v_{th}$, where the BPS prediction significantly underpredicts the observation. Experimental concerns have been ruled out as an explanation of this observation. To explain the data, it is postulated that the contribution from the ion-ion component to the total ion stopping might be significantly larger than predicted by the theories, as none of them treat both ion-ion nuclear-elastic and Coulomb scattering. A 20% increase in the total ion stopping, possibly due to ion-ion nuclear-elastic scattering, is required to explain the data, which would have an impact on our understanding of the DT-alpha energy deposition and heating of the fuel ions and would thus reduce the ignition threshold in an ignition experiment. In addition, these results indicate that the unification of the relevant physics into one theory, especially around the Bragg peak, remains challenging and an unresolved problem, even for these HED-plasma conditions. They also represent a significant advance towards providing a better understanding of DT-alpha energy deposition and heating in hot-spot ignition experiments at the NIF.

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