Controlling Quantum Confinement and Magnetic Doping of Cesium Lead Halide Perovskite Nanocrystals

Yitong Dong, David Parobek, and Dong Hee Son†

Department of Chemistry, Texas A&M University, College Station, Texas 77843, USA
(Received August 27, 2018; Accepted September 15, 2018)

ABSTRACT

Cesium lead halide (CsPbX₃) nanocrystals have emerged as a new family of semiconductor nanomaterials that can outperform existing semiconductor nanocrystals owing to their superb optical and charge transport properties. Although these materials are expected to have many superior properties, control of the quantum confinement and isoelectronic magnetic doping, which can greatly enhance their optical, electronic, and magnetic properties, has faced significant challenges. These obstacles have hindered full utilization of the benefits that can be obtained by using CsPbX₃ nanocrystals exhibiting strong quantum confinement or coupling between exciton and magnetic dopants, which have been extensively explored in many other semiconductor quantum dots. Here, we review progress made during the past several years in tackling the issues of introducing controllable quantum confinement and doping of Mn²⁺ ions as the prototypical magnetic dopant in colloidal CsPbX₃ nanocrystals.

Key words : Lead halide perovskite, Quantum dot, Magnetic doping

1. Introduction

Semiconducting perovskite nanocrystals (NCs) have attracted much attention in recent years owing to their superb light-emitting properties and facile chemical tunability of the bandgap, making them particularly suitable for applications where they act as a photon[1,2] or charge carrier source.[3,4] Among various perovskite NCs, those in the cesium lead halide (CsPbX₃) system are among the most extensively explored prototypical perovskites owing to their relative ease of synthesis and potential for outperforming more commonly used II-VI or IV-VI semiconductor NCs in various applications. In particular, their ability to harvest photons to create charge carriers, transport them long distances,[5] and emit photons[6] are superior to those of many known semiconductor NCs. Fig. 1 shows the crystal structure of CsPbX₃, transmission electron microscopy (TEM) images of colloidal NCs of CsPbX₃, and a photograph showing the color-tunable fluorescence obtained by varying the halide composition.[9] Considering the successful application of various colloidal semiconductor NCs in recent decades in solar cells,[10,11] optoelectronics,[12,13] photonics, and biological applications,[14,15] one can easily see the potential of this new family of semiconductor NCs.

Here we will focus on colloidal NCs of CsPbX₃, with particular emphasis on controlling the quantum confinement and introducing a magnetic dopant into the host NC lattice. Without question, quantum confinement of the exciton is
the most extensively explored phenomenon in various semiconductor NCs to obtain new properties unique to nanoscale semiconductor materials as distinct from their bulk counterparts. Furthermore, isoelectronic doping of paramagnetic ions into semiconductor NCs provides access to new optical, electronic, and magnetic properties not present in typical semiconductor NCs through exchange coupling between the exciton and dopant. However, introducing controlled quantum confinement and doping paramagnetic transition metal ions into CsPbX₃ NCs have been significantly more challenging than in many other semiconductor materials. Only very recently has meaningful progress been made to overcome the obstacles in performing these important tasks, enabling full exploration of the potential of this new type of semiconductor NC. Here, we review recent progress in the synthesis of colloidal quantum dots (QDs) of CsPbX₃ and Mn doping of various CsPbX₃ nanostructures.

2. Strongly Quantum-Confined CsPbX₃ NCs

In the typical hot injection synthesis condition, CsPbX₃ NCs with an edge length of > 10 nm are produced with a relatively broad size distribution. For CsPbBr₃, the exciton Bohr radius is 3.5 nm, so quantum confinement of the exciton is weak. For this reason, tuning of the bandgap in CsPbX₃ NCs was achieved primarily by varying the halide (X) composition and changing the bulk band edge level rather than varying the quantum confinement of the exciton, in contrast to the procedure for QDs with a size-tuned bandgap. The post-synthesis anion exchange reaction was often used for this purpose.

Although it is still useful to vary the absorption and emission energy in CsPbX₃ via anion exchange, other material properties that result from confinement of the exciton require a controllable NC size. For instance, the exciton binding energy, the linewidth of fluorescence emission dictated by electron-phonon coupling, and the Rashba effect prominently observed in various lead halide perovskite NCs are expected to exhibit significant dependence on the quantum confinement. To explore such properties without encountering ensemble heterogeneity and utilize them in applications, the development of a robust method for the synthesis of size-controlled NCs in the quantum-confined regime is necessary. Although post-synthesis separation, often involving tedious multistage centrifugation of the heterogeneous NC sample, can improve the ensemble homogeneity, it is a rather inefficient approach to preparing uniform perovskite QDs for practical applications.

Much effort has been made recently to extend the size range of CsPbX₃ NCs into the strongly confined regime. However, the usual kinetic control (i.e., control of the growth via the reaction time) widely utilized for the synthesis of various other semiconductor NCs has not been very effective.
effective for size control of CsPbX$_3$ NCs owing to the rapid reaction conditions under typical hot injection synthesis.\textsuperscript{9)\textsuperscript{39}} Furthermore, fast defocusing of the size after hot injection often resulted in a poor size distribution in ensembles of the NCs.\textsuperscript{\textsuperscript{39}} One strategy that has been used to gain more control over the size of the NCs in the confined regime is to lower the reaction temperature. Generally, the NC size decreased with decreasing reaction temperature, giving access to more strongly confined QDs.\textsuperscript{40–42)\textsuperscript{43}} However, the relatively large size distribution and contamination by NCs with other morphologies, such as nanoplatelets, remained an issue as reflected in the optical spectra exhibiting multiple excitonic peaks from different species (Fig. 2a). Another approach that has been taken to control the size is the use of different ligands (Fig. 2b–f).\textsuperscript{43)} The exciton absorption peak from the nanoplatelets was often prominently superimposed on the absorption from the nanocubes, whereas the nanoplatelets leave a less distinct spectroscopic signature in the luminescence spectrum owing to their relatively low emission quantum yield compared to the bright nanocubes. Recently, Manna and coworkers reported that the size of CsPbBr$_3$ perovskite nanocubes can be reduced to ca. 4 nm by varying the composition or the amount of the ligands they used, specifically, the ratio of oleic acid to oleylamine (Fig. 3a).\textsuperscript{44)\textsuperscript{45)\textsuperscript{46)\textsuperscript{47)}\textsuperscript{48)} The size was reduced by suppressing the rapid ripening rate of the perovskite NCs during growth by reducing the quantity of ligands, which the authors considered to be responsible for ripening of the NCs. Because of the difference in the rate of ripening of perovskite NCs with different halides, size control via the ligand could not be extended to CsPbX$_3$ perovskite NCs with $X = I$. Recently, Pradhan and coworkers reported that the size of CsPbBr$_3$ NCs can be reduced by adding an oleylamine–HBr mixture to the reactant. The size of the NCs could be continuously tuned by varying the amount of the oleylamine–HBr mixture (Fig. 3b).\textsuperscript{49)} They proposed that the size control mechanism is competition between Cs$^+$ and protonated ligands during growth of the NCs. All these methods rely on additional means to control the growth kinetics; therefore, the intrinsic heterogeneity of the size limits the spectroscopic homogeneity, which is crucial for various optical applications of this material.

Very recently, Son and coworkers developed a new synthesis approach that can produce strongly quantum-confined CsPbX$_3$ QDs with an unprecedented level of control over the size with ensemble uniformity.\textsuperscript{50) In contrast to the majority of other synthesis methods, which attempt to control the kinetic behavior of the reaction, the new approach developed by Son's group used the thermodynamic equilibrium of halide ions to control the size. Fig. 3 compares the optical spectra and TEM images of CsPbBr$_3$ QDs synthesized by other methods exploiting control of the kinetics.

Fig. 3. (a), (b), and (c) Absorption spectra and corresponding TEM images of QDs made in Refs. 44, 45, and 46, respectively. Images were modified with permission from Refs. 44, 45, and 46.
can be used. One crucial property of CsPbX$_3$ solids for the majority of materials. However, lead halide a viable mechanism to control the size of nanocrystalline applicability of equilibrium-based synthesis. For fixed concentrations of Cs, Pb, and X, lowering the temperature resulted in smaller QDs. This approach produces CsPbBr$_3$ QDs with high ensemble uniformity in a strongly confined regime (3.7–7 nm).

Because size control is achieved via thermodynamic equilibrium, the size of the QDs was independent of the nucleation and growth kinetics, which is one of the most distinctive and advantageous features of equilibrium-based synthesis.

**Fig. 4.** Absorption and PL spectra and TEM images of (a) CsPbCl$_3$ and (b) CsPbI$_3$ QDs obtained via halide equilibrium control method, with the corresponding TEM images. The spectra show clear excitonic features, and the TEM images show size and shape uniformity. (c) Change in size of CsPbBr$_3$ QDs with constant temperature and various Br/Pb ratios. (d) Change in size of CsPbBr$_3$ QDs with various temperatures and constant Br/Pb ratio. Images were modified with permission from Ref. 46.

In general, thermodynamic equilibrium is not considered a viable mechanism to control the size of nanocrystalline solids for the majority of materials. However, lead halide perovskite is a rare case in which the equilibrium approach can be used. One crucial property of CsPbX$_n$ NCs and related halide perovskite materials that allows the use of thermodynamic equilibrium for size control is the highly labile halide ions in the crystals. It is well known that a lead halide perovskite NC can undergo facile anion exchange even at room temperature owing to its low kinetic barrier to diffusion in the lattice. The low kinetic barrier, especially at the elevated temperatures used in hot injection (80–190°C), provides a condition in which the thermodynamic equilibrium of the halide becomes a more relevant control. The other crucial property of CsPbX$_n$ NCs is the size-dependent composition of the halide resulting from the halide-rich surface of the NCs. The relative halide content increases with decreasing QD size. This allows the equilibrium of halide ions at the interface of the NCs and the solvent medium to determine the size of the NCs. Size control was achieved either by varying the amount of halide in the reactant mixture for a given amount of Cs and Pb at a fixed reaction temperature or by varying the temperature for a fixed amount of the precursors of all three elements constituting the perovskite QDs, as shown in Fig. 4c, d, respectively.

Through the law of mass action, smaller QDs were obtained when the concentration of halides increased for a fixed reaction temperature and concentrations of Cs and Pb. For fixed concentrations of Cs, Pb, and X, lowering the temperature resulted in smaller QDs. This approach produces CsPbBr$_3$ QDs with high ensemble uniformity in a strongly confined regime (3.7–7 nm).

Because size control is achieved via thermodynamic equilibrium, the size of the QDs was independent of the nucleation and growth kinetics, which is one of the most distinctive and advantageous features of equilibrium-based synthesis. Fig. 5a, b shows the identical absorption and emission spectra of CsPbBr$_3$ QDs synthesized with various reaction quenching times and precursor injection rates, which prove that size control occurs via thermodynamic equilibrium and not kinetic control. The mechanism of size control through halide equilibrium between the QD lattice and solvent medium is illustrated by the simple model shown in Fig. 5c. In this model, control of the QD size (l) by varying the Br concentration ([Br$^{-}$]) can be viewed as satisfying the following equilibrium condition.

$$\mu_{Br^{-}_{sol}}([Br^{-}], T) = \mu_{Br^{-}_{QD}}(l, T)$$

(1)

$\mu_{Br^{-}_{sol}}([Br^{-}], T)$ is the chemical potential of Br$^{-}$ at a given [Br$^{-}$] value and temperature $T$ in the solution medium. $\mu_{Br^{-}_{QD}}(l, T)$ is the chemical potential of Br$^{-}$ in the lattice of QDs with size $l$ at temperature $T$, which depends on the relative stoichiometric composition of Br$^{-}$. $\mu_{Br^{-}_{sol}}([Br^{-}], T)$ should increase approximately logarithmically with [Br$^{-}$] following $\mu_{Br^{-}_{sol}}([Br^{-}], T) = \mu_{Br^{-}_{sol}}^{0} + RT \ln([Br^{-}])$, where

$$\mu_{Br^{-}_{sol}}^{0} = \mu_{Br^{-}_{sol}}^{0}(T)$$

(2)

$\mu_{Br^{-}_{sol}}^{0}(T)$ is the chemical potential of Br$^{-}$ in the liquid, which shows a monotonic dependence on $T$ and is expressed as

$$\mu_{Br^{-}_{sol}}^{0}(T) = \mu_{Br^{-}_{sol}}^{0}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(3)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. The equilibrium constant $K$ is expressed as

$$K = \exp \left( \frac{-\Delta G_{eq}}{RT} \right)$$

(4)

where $\Delta G_{eq}$ is the free energy change for the reaction at equilibrium, $R$ is the gas constant, and $T$ is the temperature. $K$ was determined by least squares fitting of $l$ versus $[Br^{-}]$ for each reaction temperature. $K$ is a function of $T$ and $m_{Br^{-}}$, the Br$^{-}$ concentration in the solution at equilibrium. The dependence of $K$ on $m_{Br^{-}}$ is expressed as

$$K(m_{Br^{-}}) = \frac{m_{Br^{-}}}{C_{Br^{-}}}$$

(5)

where $C_{Br^{-}}$ is the equilibrium concentration of Br$^{-}$ in solution. $K(m_{Br^{-}})$ was determined by least squares fitting of $l$ versus $[Br^{-}]$ for each reaction temperature. $K(m_{Br^{-}})$ is a function of $T$, $C_{Br^{-}}$, and $m_{Br^{-}}$.

The dependence of $m_{Br^{-}}$ on $T$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(6)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$. The dependence of $m_{Br^{-}}(T)$ on $m_{Br^{-}}(T_0)$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(7)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$. The dependence of $m_{Br^{-}}(T)$ on $m_{Br^{-}}(T_0)$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(8)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$. The dependence of $m_{Br^{-}}(T)$ on $m_{Br^{-}}(T_0)$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(9)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$. The dependence of $m_{Br^{-}}(T)$ on $m_{Br^{-}}(T_0)$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(10)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$. The dependence of $m_{Br^{-}}(T)$ on $m_{Br^{-}}(T_0)$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(11)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$. The dependence of $m_{Br^{-}}(T)$ on $m_{Br^{-}}(T_0)$ is expressed as

$$m_{Br^{-}}(T) = m_{Br^{-}}(T_0) + \sum_{i=1}^{n} \Delta H_i \ln \left( \frac{T}{T_0} \right)$$

(12)

where $\Delta H_i$ is the enthalpy of reaction for the $i$th step and $T_0$ is the reference temperature. $m_{Br^{-}}(T)$ is a function of $T$ and $m_{Br^{-}}(T_0)$.
$\mu_{\text{Br}^-}, R, \gamma$ are the chemical potential in the standard state, gas constant, and activity coefficient, respectively. The larger slope of $\mu_{\text{Br}^-,(\text{sol})}$ at lower $([\text{Br}^-])$ and the decreasing size dependence of $\mu_{\text{Br}^-,(\text{QD})}$ for larger QDs explains the less uniform size of the previously reported larger NCs synthesized by typical hot injection control of quantum confinement in Cs-based lead halide QD synthesis. Images were modified with permission from Ref. 46.

3. Mn Doping of CsPbX$_3$ NCs

Isoelectronic doping of semiconductor NCs with paramagnetic transition metal ions has been extensively investigated as a way to introduce new optical, electronic, and magnetic properties through the interaction between the exciton and d electrons of the dopant ions. For instance, II-VI QDs, such as MX (where M = Cd, Zn and X = S, Se) doped with Mn$^{2+}$ can generate intense sensitized dopant luminescence, creating a magnetically coupled exciton state enabling optical control of the magnetism, or produce energetic hot electrons via exciton-to-hot-carrier upconversion.\textsuperscript{26,49} Although it is not the focus of this article, it is worth mentioning that doping of lanthanide rare-earth metal ions has also been extensively investigated, where quantum cutting of high energy photons into low energy photons has been a topic of particular interest.\textsuperscript{49}

Because CsPbX$_3$ NCs were introduced as a new semiconductor NC material that can outperform many existing ones, doping of Mn$^{2+}$ into these materials has been extensively attempted in recent years in anticipation of the benefits seen in other semiconductor hosts. The first successful Mn doping of CsPbX$_3$ NCs was reported in CsPbCl$_3$, where hot injection synthesis of CsPbCl$_3$ NCs in the presence of an additional Mn precursor resulted in low-level doping of Mn from less than 1% up to 10%.\textsuperscript{27,29} In this synthesis, MnCl$_2$ was the most effective precursor of Mn ions, whereas many other organometallic precursors such as Mn(ac)$_2$, Mn(acac)$_2$, and Mn(oleate)$_2$ were very ineffective as dopant precursors. When MnCl$_2$ was used as the Mn precursor and CsPbBr$_3$ was used as the host material, the resulting NCs were Cl/Br mixed-halide NCs doped with Mn. However, MnBr$_2$ could not be used as a viable precursor of Mn ions for either CsPbCl$_3$ or CsPbBr$_3$ as the host under the hot injection synthesis condition, as will be discussed further later. Fig. 6a shows the absorption spectra of the undoped and Mn-doped CsPbCl$_3$ and CsPb(Cl/Br)$_3$ NCs synthesized using MnCl$_3$ as the Mn precursor at a doping concentration of less than 1%. Fig. 6b shows the corresponding luminescence spectra. In the Mn-doped NCs, characteristic Mn luminescence near 600 nm appeared, indicating successful doping of Mn into perovskite NC hosts with a sufficiently large bandgap, which can exhibit exciton–Mn energy transfer. The time-resolved Mn luminescence data shown in Fig. 6c can be fitted to a single exponential, indicating a relatively
homogeneous ligand field environment of the dopant ions, which is consistent with a low doping level. One of the issues in characterizing the detailed structure of Mn-doped perovskite NCs is quantifying the doping concentration, which can easily be overestimated owing to difficulties in removing the remaining Mn ions. Unlike many other semiconductor NCs, perovskite NCs are unstable in polar solvents, severely limiting the range of usable solvents for removing unreacted metal salts via the solvent/antisolvent approach. To overcome this issue, Son and coworkers utilized gel permeation chromatography (GPC), a technique optimized by Greytak and coworkers for cleaning II-VI QDs, to remove the excess unreacted Mn and obtain a more reliable doping concentration in the thoroughly cleaned samples. Fig. 7 shows the evolution of the absorption and emission spectra as well as the measured doping concentration after each cycle of GPC purification of the NCs. After more than four GPC cycles, the measured Mn doping concentration reached a plateau, while the same absorption and emission intensity were maintained. This demonstrates clearly the necessity for thorough removal of remaining Mn ions in the solvent for accurate quantification of the doping

![Fig. 7](image-url)

Fig. 7. (a) Absorption and (b) PL spectra of washed CsPbCl\text{Br}_{1-x} sample. (c) Tested Mn doping concentration of samples with different numbers of washings. Images were modified with permission from Ref. 27.

![Fig. 8](image-url)

Fig. 8. (a) Absorption and PL spectra of CsPbCl\text{Br}_{1-x} sample with high Mn doping concentration (~ 46%). (b) TEM images of CsPbCl\text{Br}_{1-x} NCs with high Mn doping concentration. Inset shows photograph of Mn-doped CsPbCl\text{Br}_{1-x} NCs under UV excitation. Images were modified with permission from Ref. 51. (c) Method to dope Mn in CsPbCl\text{Br}_{1-x} NCs using an ammonium chloride mixture. (d) PL spectra of CsPbCl\text{Br}_{1-x} NCs acquired with increasing amount of ammonium chloride mixture. Images were modified with permission from Ref. 52.
concentration. In later works on Mn-doped CsPbCl$_3$ NCs, the hot injection synthesis protocol was further optimized to increase the doping concentration. Yang and coworkers were able to reach a doping concentration of 46% and luminescence quantum yield of 56% by using a large amount of excess Mn precursors (e.g., a Pb/Mn ratio of 1:10). They showed the absorption and emission spectra and TEM images of the highly Mn-doped CsPbCl$_3$ NCs. Pradhan and coworkers took another approach to increasing the efficiency of Mn doping, where they used oleylammominium chloride as an additional reactant that can boost the doping efficiency. They observed a large increase in the Mn luminescence intensity with increasing amount of oleylammominium chloride for a given amount of Mn precursor. The authors attributed the increased efficiency of Mn doping to the formation of a higher concentration of NCs during the early phase of the reaction, which increased the probability of Mn doping, although the validity of this explanation is not yet fully established.

More recent synthesis of Mn-doped CsPbCl$_3$ has expanded the reaction conditions beyond the hot injection condition. Meijerink and coworkers showed examples of room-temperature synthesis methods for Mn doping of CsPbCl$_3$, with non-halide Mn precursors. In their first report, they used metal acetate salts as the precursor, which were converted to metal oleate complexes in the presence of ligands, and then added HCl to protonate the carboxylate group, increasing the amount of monomer initiating the formation of nanocubes. HCl also created a Cl-rich surface, supplying ample binding sites for Mn ions and facilitating Mn doping. Further coating of Mn-doped CsPbCl$_3$ with an additional CsPbCl$_3$ shell improved the Mn luminescence quantum yield.

Although earlier efforts to synthesize Mn-doped CsPbCl$_3$ NCs produced relatively large cube-shaped NCs that experience very weak quantum confinement, several recent studies reported the synthesis of CsPbCl$_3$ nanoplatelets with strong confinement in one dimension. Nag and coworkers synthesized Mn-doped CsPbCl$_3$ nanoplatelets with a thickness of 2.2 nm, which imposes strong confinement along the thickness direction, by a room-temperature reaction. They were also able to vary the doping concentration in the range of 0.2–2%. Pradhan and coworkers reported another method of generating Mn-doped CsPbCl$_3$ nanoplatelets, which involves the initial synthesis of a Mn-doped monolayer structure and subsequent formation of nanoplatelets by the addition of Cs-oleate. They produced 5-nm-thick nanoplatelets with a relatively large lateral dimension in the range of 20–580 nm, which varied with the concentration of Cs and Mn during the reaction. Because quantum confinement of the exciton in Mn-doped semiconductor NCs can enhance the exciton–donor interaction, which determines various magneto-optical properties, continued progress in the synthesis of strongly confined Mn-doped perovskite NCs will be important for expanding their applicability.

Most of the work on Mn doping of CsPbBr$_3$ NCs has focused on CsPbCl$_3$, despite its less favorable optical properties, such as a higher bandgap and lower luminescence quantum yield, which are presumably due to its higher defect density compared to other halide perovskites. This is because doping of Mn is most favorable in CsPbCl$_3$ and becomes increasingly difficult for bromide and iodide perovskite hosts. To date, successful doping of Mn in a CsPbBr$_3$ NC host could not be achieved by simply extending the doping method used for producing Mn-doped CsPbCl$_3$ NCs. Klimov and coworkers hypothesized that direct hot injection synthesis of Mn-doped CsPbBr$_3$ using MnBr$_2$ was energetically unfavorable owing to the large difference in bond energy between Pb–Br (249 kJ/mol) and Mn–Br (314 kJ/mol) compared to that between Pb–Cl (301 kJ/mol) and Mn–Cl (338 kJ/mol). They argued that the higher stability of the Mn–Br bond compared to the Pb–Br bond impeded the incorporation of Mn$^{2+}$ into the CsPbBr$_3$ lattice. Although it is still not clearly understood yet, the difference in the thermochemistry can certainly be a factor affecting the ability to dope Mn$^{2+}$ into different CsPbX$_3$ NCs.

There are several reports of attempts to utilize the ion exchange reaction to obtain Mn-doped CsPbBr$_3$ NCs. For instance, anion exchange of Mn-doped CsPbCl$_3$ with Br$^-$ resulted in not only the exchange of anions but also removal of Mn ions in the host lattice. Others attempted cation exchange of Pb with Mn in CsPbBr$_3$ using MnCl$_4$. However, this also resulted in anion exchange, ultimately forming Mn-doped NCs with Cl$^-$ as the dominant anion in the lattice. Chen and coworkers claimed that they were able to dope Mn into CsPbBr$_3$ and that it increased the stability of the NCs significantly, although they did not observe Mn emission from their sample.

Recently, Son and coworkers developed a more robust synthesis method that produces Mn-doped CsPbBr$_3$ NCs that exhibit intense Mn luminescence. The synthesis was a two-step procedure in which they first synthesized a Mn-doped monolayer structure ($L_4$[Pb$_{1-x}$Mn$_x$Br$_4$], where L is a ligand) as an intermediate, which they converted to Mn-doped CsPbBr$_3$ NCs by adding Cs-oleate (Fig. 9a). In this method, HBr was used to facilitate the formation of a stable $L_4$[Pb$_{1-x}$Mn$_x$Br$_4$] monolayer structure. Mn doping of this monolayer structure was clear from the observation of intense Mn luminescence upon excitation of the host structure at 370 nm. The absorption and emission spectra of $L_4$[Pb$_{1-x}$Mn$_x$Br$_4$] are shown in Fig. 9b, c, respectively. The absorption spectrum, which has a sharp absorption peak at 395 nm, is very similar to that of a $L_4$[PbBr$_4$] monolayer.
reported previously, indicating that the $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ is essentially a Mn-doped monolayer structure of $L_2[\text{PbBr}_4]$. On the other hand, the luminescence is centered at 620 nm and arises entirely from Mn. The similarity of the absorption and photoluminescence excitation (PLE) spectra of $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ proved that the Mn luminescence in this monolayer structure is also sensitized. The intensity of the Mn luminescence in the $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ monolayer structure was well correlated with the amount of HBr used, which indicated that HBr facilitates doping of Mn into this structure. $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ often forms a multilayer structure when dried, as indicated by a series of small-angle diffraction peaks separated by $2\theta = 2.2^\circ$ in the powder X-ray diffraction patterns, which corresponds to a 4.1 nm spacing between the stacked layers separated by the ligand. When Cs-oleate was added to the solution of the $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ monolayer structure at 200°C, Mn-doped CsPbBr$_3$ NCs were formed. Interestingly, the resulting product was a mixture of Mn-doped CsPbBr$_3$ NCs with two different morphologies, i.e., nanocubes (6.5–8.5 nm) and nanoplatelets (~2 nm in thickness). However, the two species could be readily separated via centrifugation, providing access to NCs of two different dimensions. The Mn luminescence intensity of the Mn-doped CsPbBr$_3$ nanocubes was well correlated with that of the $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ monolayer structure, showing an increase with the amount of HBr used during the initial formation of the $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ monolayer structure (Fig. 9d–f). This suggests that the formation of $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$, which contains the same Mn–Br coordination present in Mn-doped CsPbBr$_3$ NCs, is crucial to formation of Mn-doped CsPbBr$_3$ NCs under a hot injection condition. The robustness of the Mn–Br bond in $L_2[\text{Pb}_x\text{Mn}_{1-x}\text{Br}_4]$ was demonstrated by the strong Mn luminescence, which survived even at 120°C. It is also notable that unlike the Mn-doped CsPbBr$_3$ nanocubes, which exhibited both exciton and Mn luminescence, the Mn-doped CsPbBr$_3$ nanoplatelets exhibited primarily Mn luminescence, as can be determined easily from the color of

Fig. 9. (a) Schematic of Mn doping of CsPbBr$_3$ NCs. (b) Absorption and (c) PL spectra of Mn-doped CsPbBr$_3$ NCs acquired with different amounts of HBr. (d) Time-dependent PL intensity of Mn phosphorescence of samples obtained with different amounts of HBr. Images were modified with permission from Ref. 28.
the luminescence (Fig. 9a). This may reflect the potentially more rapid exciton–Mn energy transfer rate in the partially quantum confined nanoplatelets compared to that in the only weakly quantum confined nanocubes.

Because doping of Mn$^{2+}$ ions into semiconductor NCs creates new optical, electronic, and magnetic properties resulting from exchange coupling between the exciton and dopant, the magnitude of the coupling is an important subject of investigation. It was explored quite extensively in II-VI Mn-doped semiconductor NCs as a function of the structural parameters, which affect the spatial overlap of the exciton and dopant wavefunction dictating the strength of the exchange coupling of the exciton and dopant.\textsuperscript{60,61} For the Mn-doped CsPbX$_3$ NC system, Son and coworkers performed the first experiment to provide useful information on the strength of the exciton–dopant exchange coupling in comparison to that in Mn-doped II-VI semiconductor NCs.\textsuperscript{62} The approach taken was to compare the exciton–Mn energy transfer times ($\tau_{ET}$) of Mn-doped CsPbCl$_3$ NCs and Mn-doped CdS/ZnS NCs after normalizing to the doping concentration. In their study, $\tau_{ET}$ was obtained by comparative analysis of the exciton relaxation dynamics in Mn-doped and undoped NCs using pump–probe transient absorption spectroscopy to monitor the dynamics of exciton relaxation. The energy transfer pathway that exists only in Mn-doped NCs was manifested as an additional dynamic component resulting from exchange coupling between the exciton and magnetic dopant ions that are not present in undoped perovskite NCs. Although the ability to control the quantum confinement and introduce magnetic dopants into the lattice of CsPbX$_3$ NCs was acquired only recently, methodologies have rapidly expanded to CsPbX$_3$ NCs of different dimensions beyond simple nanocubes. With continuing development of the synthesis of CsPbX$_3$ NCs with higher levels of structural and chemical control, it will also be necessary to explore the structure–property relationship and use the new properties in applications.

4. Outlook and Conclusion

Although the first chemical synthesis of colloidal CsPbX$_3$ perovskite NCs was reported only several years ago, much progress has been made since then in both the synthesis and characterization of CsPbX$_3$ NCs with much higher control of the dimensions and chemical composition. One of the most important achievements in the synthesis of CsPbX$_3$ perovskite NCs was the development of methods to precisely control the quantum confinement with very high ensemble uniformity of the morphology and size in the strongly confined regime. This opened the door to exploitation of quantum confinement of excitons in perovskite NCs for their application, which had been lagging compared to that of other semiconductor NCs. Another notable achievement in the synthesis of colloidal CsPbX$_3$ perovskite NCs was successful doping with paramagnetic transition metal ions such as Mn$^{2+}$. This provided access to various new optical, electronic, and magnetic properties originating from the coupling of the exciton and magnetic dopant ions that are not present in undoped perovskite NCs. Although the ability to control the quantum confinement and introduce magnetic dopants into the lattice of CsPbX$_3$ NCs was acquired only recently, methodologies have rapidly expanded to CsPbX$_3$ NCs of different dimensions beyond simple nanocubes.

Acknowledgments

This work was supported by the Welch Foundation (A-1639). D.P. was supported by the NSF GRFP (DGE-1252521).

REFERENCES

1. X. Zhang, H. Lin, H. Huang, C. Reckmeier, Y. Zhang, W. C. H. Choy, and A. L. Rogach, “Enhancing the Brightness of Cesium Lead Halide Perovskite Nanocrystal Based Green Light-Emitting Devices through the Interface Engineering with Perfluorinated Ionomer,” \textit{Nano Lett.}, \textbf{16} [2] 1415–20 (2016).

2. G. Li, F. W. R. Rivarola, N. J. L. K. Davis, S. Bai, T. C. Jellicoe, F. de la Peña, S. Hou, C. Ducati, F. Gao, R. H. Friend, N. C. Greenham, and Z.-K. Tan, “Highly Efficient Perovskite Nanocrystal Light-Emitting Diodes Enabled by a Universal Crosslinking Method,” \textit{Adv. Mater.}, \textbf{28} [18] 3528–34 (2016).

Fig. 10. Scheme of exciton-to-dopant energy transfer in Mn-doped CsPbCl$_3$ NCs and transient absorption data for doped and undoped CsPbCl$_3$ NCs. Images were modified with permission from Ref. 62.
3. Q. A. Akkerman, M. Gandini, F. Di Stasio, P. Rastogi, F. Palazon, G. Bertoni, J. M. Ball, M. Prato, A. Petrozza, and L. Manna, “Strongly Emissive Perovskite Nanocrystal Inks for High-Voltage Solar Cells,” *Nat. Energy* **2** 16194 (2016).

4. A. Swarnkar, A. R. Marshall, E. M. Sanehira, B. D. Cher nomordik, and D. T. Moore, “Quantum Dot–Induced Phase Stabilization of α-CsPbI$_3$ Perovskite for High-Efficiency Photovoltaics,” *Science*, **354** [6308] 92–5 (2016).

5. A. Kojima, K. Teshima, Y. Shirai, and T. Miyasaka, “Organometal Halide Perovskites as Visible-Light Sensitizers for Photovoltaic Cells,” *J. Am. Chem. Soc.*, **131** [17] 6050–51 (2009).

6. G. R. Yettapu, D. Talukdar, S. Sarkar, A. Swarnkar, A. Nag, P. Ghosh, and P. Mandal, “Terahertz Conductivity within Colloidal CsPbBr$_3$ Perovskite Nanocrystals: Remarkably High Carrier Mobilities and Large Diffusion Lengths,” *Nano Lett.*, **16** [8] 4838–48 (2016).

7. S. ten Brinck and I. Infante, “Surface Termination, Morphology, and Bright Photoluminescence of Cesium Lead Halide Perovskite Nanocrystals,” *ACS Energy Lett.*, **1** [6] 1266–72 (2016).

8. R. E. Brandt, J. R. Poindexter, P. Gorai, R. C. Kurchin, R. L. Z. Hoye, L. Nienhaus, M. W. B. Wilson, J. A. Polizotto, R. Sereika, R. Zaltskaukas, L. C. Lee, J. L. MacManus-Driscoll, M. Bawendi, V. Stevanović, and T. Buonassisi, “Searching for “Defect-Tolerant” Photovoltaic Materials: Combined Theoretical and Experimental Screening,” *Chem. Mater.*, **29** [11] 4667–74 (2017).

9. L. Protescu, S. Yakunin, M. I. Bodnarchuk, F. Krieg, R. Caputo, C. H. Hendon, R. X. Yang, A. Walsh, and M. V. Kovalenko, “Nanocrystals of Cesium Lead Halide Perovskites (CsPbX$_3$, X = Cl, Br, and I): Novel Optoelectronic Materials Showing Bright Emission with Wide Color Gamut,” *Nano Lett.*, **15** [6] 3692–96 (2015).

10. P. K. Santra and P. V. Kamat, “Mn-Doped Quantum Dot Sensitized Solar Cells: A Strategy to Boost Efficiency over 5%,” *J. Am. Chem. Soc.*, **134** [5] 2508–11 (2012).

11. P. V. Kamat, “Quantum Dot Solar Cells. The Next Big Thing in Photovoltaics,” *J. Phys. Chem. Lett.*, **4** [6] 908–18 (2013).

12. A. J. Nozik, “Quantum Dot Solar Cells,” *Phys. E,* **14** [1–2] 115–20 (2002).

13. A. H. Ip, S. M. Thon, S. Hoogland, O. Voznyy, D. Zhitomirsky, R. Debnath, L. Levina, L. R. Rollny, G. H. Carey, A. Fischer, K. W. Kemp, I. J. Kramer, Z. Ning, A. J. Labelle, W. K. Chou, A. Amassian, and E. H. Sargent, “Hybrid Passivated Colloidal Quantum Dot Solids,” *Nat. Nanotechnol.*, **7** 577 (2012).

14. I. Gur, N. A. Fromer, M. L. Geier, and A. P. Alivisatos, “Air-Stable All-Inorganic Nanocrystal Solar Cells Processed from Solution,” *Science*, **310** [5747] 462–5 (2005).

15. J.-H. Choi, H. Wang, S. J. Oh, T. Paik, P. Sung, J. Sung, X. Ye, T. Zhao, B. T. Diroll, C. B. Murray, C. R. Kagan, “Exploiting the Colloidal Nanocrystal Library to Construct Electronic Devices,” *Science*, **352** [6282] 205–8 (2016).

16. C. R. Kagan, E. Lifshitz, E. H. Sargent, and D. V. Talapin, “Building Devices from Colloidal Quantum Dots,” *Science*, **353** [6302] aac5523 (2016).

17. S. A. McDonald, G. Konstantatos, S. Zhang, P. W. Cyr, E. J. D. Klem, L. Levina, E. H. Sargent, “Solution-Processed PbS Quantum Dot Infrared Photodetectors and Photovoltaics,” *Nat. Mater.*, **4** 138 (2005).

18. X. Michalet, F. F. Pinaud, L. A. Bentolila, J. M. Tsay, S. Doose, J. J. Li, G. Sundaresan, A. M. Wu, S. S. Gambhir, and S. Weiss, “Quantum Dots for Live Cells, in vivo Imaging, and Diagnostics,” *Science* **307** [5709] 538–44 (2005).

19. I. L. Medintz, H. T. Uyeda, E. R. Goldman, and H. Mato ussi, “Quantum Dot Bioconjugates for Imaging, Labelling and Sensing,” *Nat. Mater.*, **4** [6] 435–46 (2005).

20. X. H. Gao, Y. Y. Cui, R. M. Levenson, L. W. K. Chung, and S. M. Nie, “In vivo Cancer Targeting and Imaging with Semiconductor Quantum Dots,” *Nat. Biotechnol.*, **22** [8] 969–76 (2004).

21. A. P. Alivisatos, “Semiconductor Clusters, Nanocrystals, and Quantum Dots,” *Science*, **271** [5251] 933–37 (1996).

22. D. J. Norris, N. Yao, F. T. Charnock, and T. A. Kennedy, “High-Quality Manganese-Doped ZnSe Nanocrystals,” *Nano Lett.*, **1** [1] 3–7 (2001).

23. D. J. Norris, A. L. Efros, and S. C. Erwin, “Doped Nanocrystals,” *Science*, **319** [5871] 1776–79 (2008).

24. N. Pradhan and X. G. Peng, “Efficient and Color-Tunable Mn-Doped ZnSe Nanocrystal Emitters: Control of Optical Performance via Greener Synthetic Chemistry,” *J. Am. Chem. Soc.*, **129** [11] 3339–47 (2007).

25. B. B. Srivastava, S. Jana, N. S. Karan, S. Paria, N. R. Jana, D. S. Sarma, and N. Pradhan, “Highly Luminescent Mn-Doped ZnS Nanocrystals: Gram-Scale Synthesis,” *J. Phys. Chem. Lett.*, **1** [9] 1454–58 (2010).

26. R. Beaulac, L. Schneider, P. I. Archer, G. Bacher, and D. R. Gamelin, “Light-Induced Spontaneous Magnetization in Doped Colloidal Quantum Dots,” *Science*, **325** [5943] 973–76 (2009).

27. D. Parobek, B. B. Roman, Y. Dong, H. Jin, E. Lee, M. Sheldon, and D. H. Son, “Exciton-to-Dopant Energy Transfer in Mn-Doped Cesium Lead Halide Perovskite Nanocrystals,” *Nano Lett.*, **16** [12] 7376–80 (2016).

28. D. Parobek, Y. Dong, T. Qiao, and D. H. Son, “Direct Hot-Injection Synthesis of Mn-Doped CsPbBr$_2$ Nanocrystals,” *Chem. Mater.*, **30** [9] 2939–44 (2018).

29. W. Liu, Q. Lin, H. Li, K. Wu, I. Robel, J. M. Pieryga, and V. I. Klimov, “Mn$^2+$-Doped Lead Halide Perovskite Nanocrystals with Dual-Color Emission Controlled by Halide Content,” *J. Am. Chem. Soc.*, **138** [45] 14954–61 (2016).

30. Q. A. Akkerman, V. D'Innocenzo, S. Accornero, A. Scar pelini, A. Petrozza, M. Prato, and L. Manna, “Tuning the Optical Properties of Cesium Lead Halide Perovskite Nanocrystals by Anion Exchange Reactions,” *J. Am. Chem. Soc.*, **137** [32] 10276–81 (2015).

31. G. Nedelev, L. Protescu, S. Yakunin, M. I. Bodnarchuk, M. J. Grotevent, and M. V. Kovalenko, “Fast Anion Exchange in Highly Luminescent Nanocrystals of Cesium Lead Halide Perovskites (CsPbX$_3$, X = Cl, Br, I),” *Nano Lett.*, **15** [8] 5635–40 (2015).

32. J. Li, L. Luo, H. Huang, C. Ma, Z. Ye, J. Zeng, and H. He, “2D Behaviors of Excitons in Cesium Lead Halide Perovskite Nanoplatelets,” *J. Phys. Chem. Lett.*, **8** [6] 1161–
33. A. Pan, B. He, X. Fan, Z. Liu, J. J. Urban, A. P. Alivisatos, L. He, and Y. Liu, “Insight into the Ligand-Mediated Synthesis of Colloidal CsPbBr3 Nanoplatelets,” ACS Nano, 10 [8] 7943–54 (2016).
34. S. Sun, D. Yuan, Y. Xu, A. Wang, and Z. Deng, “Ligand-Mediated Synthesis of Shape-Controlled Cesium Lead Halide Perovskite Nanocrystals via Reprecipitation Process at Room Temperature,” ACS Nano, 10 [3] 3648–57 (2016).
35. G. Almeida, L. Goldoni, Q. Akkerman, Z. Dang, A. H. Khan, S. Marras, I. Moreels, L. Manna, “Role of Acid–Base Equilibria in the Size, Shape, and Phase Control of Cesium Lead Halide Perovskite Nanocrystals,” ACS Nano, 12 [2] 1704–11 (2018).
36. A. Dutta, S. K. Dutta, S. Das Adhikari, and N. Pradhan, “Tuning the Size of CsPbBr3 Nanocrystals: All at One Constant Temperature,” ACS Energy Lett., 3 [2] 329–34 (2018).
37. Y. Dong, T. Qiao, D. Kim, D. Parobek, D. Rossi, and D. H. Son, “Precise Control of Quantum Confinement in Cesium Lead Halide Perovskite Quantum Dots via Thermodynamic Equilibrium,” Nano Lett., 18 [6] 3716–22 (2018).
38. D. Parobek, Y. Dong, T. Qiao, D. Rossi, and D. H. Son, “Photoinduced Anion Exchange in Cesium Lead Halide Perovskite Nanocrystals,” J. Am. Chem. Soc., 139 [12] 4358–61 (2017).
39. Y. Dong, J. Choi, H.-K. Jeong, and D. H. Son, “Hot Electrons Generated from Doped Quantum Dots via Upconversion of Excitons to Hot Charge Carriers for Enhanced Photocatalysis,” J. Am. Chem. Soc., 137 [16] 5549–54 (2015).
40. G. Pan, X. Bai, D. Yang, X. Chen, P. Jing, S. Qu, L. Zhang, D. Zhou, J. Zhu, W. Xu, B. Dong, and H. Song, “Doping Lanthanide into Perovskite Nanocrystals: Highly Improved and Expanded Optical Properties,” Nano Lett., 17 [12] 8005–11 (2017).
41. Y. Shen, M. Y. Gee, R. Tan, P. J. Pellechia, and A. B. Greytak, “Purification of Quantum Dots by Gel Permeation Chromatography and the Effect of Excess Ligands on Shell Growth and Ligand Exchange,” Chem. Mater., 25 [14] 2838–48 (2013).
42. H. Liu, Z. Wu, J. Shao, D. Yao, H. Gao, Y. Liu, W. Yu, H. Zhang, and B. Yang, “CsPbMnxClyCl1-y Perovskite Quantum Dots with High Mn Substitution Ratio,” ACS Nano, 11 [2] 2329–47 (2017).
43. S. Das Adhikari, S. K. Dutta, A. Dutta, A. K. Guria, and N. Pradhan, “Chemically Tailoring the Dopant Emission in Manganese-Doped CsPbCl3 Perovskite Nanocrystals,” Angew. Chem. Int. Ed., 56 [30] 8746–50 (2017).
44. K. Xu, C. C. Lin, X. Xie, and A. Meijerink, “Efficient and Stable Luminescence from Mn2+ in Core and Core-Isocristalline Shell CsPbCl3 Perovskite Nanocrystals,” Chem. Mater., 29 [10] 4265–72 (2017).
45. W. J. Mir, M. Jagadeeswararao, S. Das, and A. Nag, “Colloidal Mn-Doped Cesium Lead Halide Perovskite Nanoplatelets,” ACS Energy Lett., 2 [8] 537–43 (2017).
46. S. Das Adhikari, A. Dutta, S. K. Dutta, and N. Pradhan, “Layered Perovskites Lx4(Pb1-xMnx)4Cl6 to Mn-Doped CsPbCl3 Perovskite Platelets,” ACS Energy Lett., 3 [6] 1247–53 (2018).
47. F. Li, Z. Xia, C. Pan, Y. Gong, L. Gu, Q. Liu, and J. Z. Zhang, “High Br Content CsPb(Cl,Br)x Nanoparticles with Strong Mn2+ Emission through Diverse Cation/Anion Exchange Engineering,” ACS Appl. Mater. Interfaces, 10 [14] 11739–46 (2018).
48. S. Zhu, Y. Liu, J. Li, C. Liu, R. Feng, F. Jiang, Y. Li, J. Song, H. Zeng, M. Hong, and X. Chen, “Stabilizing Cesium Lead Halide Perovskite Lattice through Mn(II) Substitution for Air-Stable Light-Emitting Diodes,” J. Am. Chem. Soc., 139 [33] 11443–50 (2017).
49. K. Xu and A. Meijerink, “Tuning Exciton–Mn2+ Energy ...
Transfer in Mixed Halide Perovskite Nanocrystals,” Chem. Mater., 30 [15] 5346–52 (2018).
59. R. Beaulac, P. I. Archer, J. van Rijssel, A. Meijerink, and D. R. Gamelin, “Exciton Storage by Mn²⁺ in Colloidal Mn²⁺-Doped CdSe Quantum Dots,” Nano Lett., 8 [9] 2949–53 (2008).
60. H.-Y. Chen, S. Maiti, and D. H. Son, “Doping Location-Dependent Energy Transfer Dynamics in Mn-Doped CdS/ZnS Nanocrystals,” ACS Nano, 6 [1] 583–91 (2012).
61. W. D. Rice, W. Liu, V. Pinchetti, D. R. Yakovlev, V. I. Klimov, and S. A. Crooker, “Direct Measurements of Magnetic Polarons in Cd₁₋ₓMnₓSe Nanocrystals from Resonant Photoluminescence,” Nano Lett., 17 [5] 3068–75 (2017).
62. D. Rossi, D. Parobek, Y. Dong, and D. H. Son, “Dynamics of Exciton–Mn Energy Transfer in Mn-Doped CsPbCl₃ Perovskite Nanocrystals,” J. Phys. Chem. C, 121 [32] 17143–49 (2017).