Inverse correlation between quasiparticle mass and $T_c$ in a cuprate high-$T_c$ superconductor

Carsten Putzke,1 Liam Malone,1 Sven Badoux,2 Baptiste Vignolle,2 David Vignolles,2 Wojciech Tabis,2,3 Philip Walmsley,1 Matthew Bird,1 Nigel E. Hussey,4 Cyril Proust,2 Antony Carrington1*

Close to a zero-temperature transition between ordered and disordered electronic phases, quantum fluctuations can lead to a strong enhancement of electron mass and to the emergence of competing phases such as superconductivity. A correlation between the existence of such a quantum phase transition and superconductivity is quite well established in some heavy fermion and iron-based superconductors, and there have been suggestions that high-temperature superconductivity in copper-oxide materials (cuprates) may also be driven by the same mechanism. Close to optimal doping, where the superconducting transition temperature $T_c$ is maximal in cuprates, two different phases are known to compete with superconductivity: a poorly understood pseudogap phase and a charge-ordered phase. Recent experiments have shown a strong increase in quasiparticle mass at optimal doping, suggesting that quantum fluctuations of the charge-ordered phase may be responsible for the high-$T_c$ superconductivity. We have tested the robustness of this correlation by performing quantum oscillation studies on the stoichiometric compound YBa$_2$Cu$_4$O$_8$ under hydrostatic pressure. In contrast to the results for YBa$_2$Cu$_3$O$_{7-x}$, we find that in YBa$_2$Cu$_4$O$_8$ the mass decreases as $T_c$ increases under pressure. This inverse correlation between $m^*$ and $T_c$ suggests that quantum fluctuations of the charge order enhance $m^*$ but do not enhance $T_c$.

INTRODUCTION

In a variety of systems such as organics, heavy fermions, iron pnictides, and copper oxides (cuprates), superconductivity is found in close proximity to an antiferromagnetic phase (1–4). Starting with the non-superconducting parent antiferromagnetic phase, because the material properties are tuned by the application of pressure or chemical doping, the magnetic ordering temperature $T_N$ decreases. The doping pressure where $T_N$ extrapolates to zero is known as a quantum critical point (QCP). In many systems, the superconducting transition temperature $T_c$ is maximal close to the doping pressure where the QCP is located, thus suggesting an intimate connection between the two. Close to a QCP, quantum fluctuations between the ordered magnetic phase and the disordered paramagnetic phase become strong, and if such fluctuations produce electron pairing, this may enhance $T_c$. The strong fluctuations also affect the normal-state properties, causing a strong increase in quasiparticle mass and in the non-Fermi liquid behavior of the transport properties. Hence, close to a QCP, the normal-state entropy is increased, and this too could provide a boost for $T_c$ (1, 3).

The applicability of such a quantum critical scenario to cuprate superconductors is somewhat debatable because in these materials, $T_c$ is zero close to the QCP of the antiferromagnetic phase. Instead, $T_c$ is maximal close to the end points of two other phases: the pseudogap phase and the charge density wave (CDW) phase. Whether there is a QCP at the end points of either of these phases is unclear, as is their exact position (5, 6). The nature and origin of the long-established pseudogap phase, which can be characterized as a depression of the density of states at the Fermi level (7) and a decoherence of quasiparticles on certain parts of the Fermi level (8), are not well understood. In contrast, the more recently discovered CDW phases (9–11) seem to be more conventional, although many aspects, such as the microscopic mechanism of their formation and their variation with pressure, still need to be clarified. There are two distinct CDW phases. One phase, which is essentially two-dimensional, forms independently of a magnetic field below $T \sim 15$ K (10, 11), whereas a second distinct three-dimensional order forms for fields above $T \sim 15$ T and below $T \sim 60$ K (9, 12, 13).

A CDW produces an additional periodic potential, which reduces the size of the Brillouin zone and folds the Fermi surface into small Fermi pockets. There is compelling evidence that a CDW-induced Fermi surface reconstruction is responsible for the small Fermi surface pockets, which are observed in quantum oscillation (QO) measurements (14), so QOs give a precise way of characterizing how the normal-state electronic structure is affected by the CDW. In particular, from the thermal damping of the QO amplitude, the quasiparticle mass $m^*$ can be inferred. $m^*$ is directly related to the renormalization of the band structure close to the Fermi level coming from the interaction of the electrons with collective excitations such as phonons or spin/CDW fluctuations.

Recently, the variation of $m^*$ in the cuprate YBa$_2$Cu$_4$O$_{7-\delta}$ (Y123) was measured over a wide range of hole doping ($n_p$) (15, 16). A dome-shaped variation of $(m^*_p)^{-1}$ versus $n_p$ was found, with $(m^*_p)^{-1}$ extrapolating to zero at two points, $n_p \approx 0.08$ and $n_p \approx 0.18$. This is suggestive of QCPs close to $n_p \approx 0.165$ and $n_p \approx 0.08$ and because the latter value is close to the value ($n_p \approx 0.165$) where $T_c$ is maximal (17), this may indicate that the interactions responsible for the mass enhancement also give rise to the high-$T_c$ superconductivity (18). It is natural to assume that these interactions are the fluctuations of the CDW, although additional coupling to the pseudogap fluctuations is also a possible cause (15).

1H. H. Wills Physics Laboratory, University of Bristol, Tyndall Avenue, Bristol BS8 1TL, UK.
2Laboratoire National des Champs Magnétiques Intenses, CNRS-INSA-UJF-UPS, 31400 Toulouse, France. 3AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, aleja Adama Mickiewicza 30, 30-059 Krakow, Poland. 4High Field Magnet Laboratory (HFML-EMFL), Radboud University, Toernooiveld 7, 6525 ED Nijmegen, Netherlands.
5Corresponding author. E-mail: a.carrington@bristol.ac.uk
Although chemical doping (for example, by changing the oxygen content \( \delta \) in Y123) is the most widely used method of tuning the properties of cuprates, pressure is another available tuning parameter, and indeed the highest transition temperatures are achieved under high pressure (19). Here, we report measurements of QOs under high pressure in the stoichiometric cuprate YBa\(_2\)Cu\(_4\)O\(_8\) (Y124). This allows us to study how the electronic correlations evolve as \( T_c \) is increased with pressure in a single sample with fixed stoichiometry.

Y124 has a structure similar to that of Y123 but with a double copper oxide chain layer. Unlike Y123, in Y124 the oxygen stoichiometry is fixed and thus cannot be used to change the doping or \( T_c \). However, the application of hydrostatic pressure to Y124 leads to a strong increase in \( T_c \) by 5.5 K/GPa at low pressure (20). \( T_c \) increases from 79 K at ambient pressure [where \( n_p \approx 0.13(1) \) (21)] to \( \approx 105 \) K at \( \approx 8 \) GPa (22, 23) at which point it undergoes a structural phase transition and becomes nonsuperconducting (24, 25). Although direct x-ray or nuclear magnetic resonance (NMR) evidence for the existence of a CDW in Y124 remains missing to date, its existence is likely, as the transport properties (26) show features similar to those in Y123, which have been ascribed to CDW formation. Furthermore, QOs with a frequency similar to those seen in Y123 have been observed in Y124 (21, 27), strongly suggesting that a similar CDW-driven Fermi surface reconstruction occurs.

**RESULTS**

Figure 1 shows ambient-pressure measurements of the \( c \)-axis resistivity \( \rho_c \) in a pulsed magnetic field of up to 67 T applied parallel to the \( c \) axis (perpendicular to the CuO\(_2\) planes) of a single-crystal sample of Y124. Shubnikov–de Haas QOs in the resistance are clearly visible in the raw data above \( \approx 52 \) T but are more clearly seen in \( d\rho_c/dH \) (Fig. 1, inset). The frequency \( F \) (in the inverse field) of these oscillations is 640 T, consistent with previous reports (21, 27). The size of the oscillations is remarkably reproducible between different crystals (fig. S1). By fitting the temperature dependence of the oscillation amplitude, the quasiparticle mass can be extracted (Fig. 1B). We find a value of \( m^* = 1.80(5) \) (in units of free electron mass). A value of \( m^* = 1.9(1) \) was measured in two further samples (fig. S2). This is consistent with another recent study (28); however, older measurements (21, 27) found larger values of \( m^* \), but the errors were considerably larger than in the present work (see also Materials and Methods).

This same sample was then transferred to a nonmetallic pressure cell (29). Measurements of \( \rho_c(H,T) \) at different applied pressures \( p \) up to \( p = 0.84 \) GPa are shown in Fig. 2. Because of the larger bore of the magnet used for the pressure cell measurements, the maximum field here was 58 T. At room temperature, \( \rho_c \) decreases linearly with increasing pressure (fig. S3A), and this decrease is essentially independent of temperature down to \( T_c(H = 0) \) (fig. S3B). However, when superconductivity is suppressed by field at low temperatures, a markedly different behavior emerges. As can be seen in Fig. 2, both the magnitude of the resistivity and its slope \( (d\rho_c/dH) \) at 2.5 K and 50 T almost double between \( p = 0 \) and \( p = 0.37 \) GPa (the lowest pressure that we are able to apply), which suggests that, at low temperatures, the size of magnetoresistance significantly increases with the application of pressure, presumably due to an increase in the mean free path on some parts of the Fermi surface. This result was reproduced in three samples.

We were able to observe QOs at all pressures. The oscillatory part of the magnetoresistance at a fixed temperature of 2.5 K is shown in
Fig. 3A. The relative oscillation amplitude ($\Delta \rho_c/\rho_c$) increases by around a factor of 2 between $p = 0$ and $p = 0.37$ GPa but then remains approximately constant as pressure is increased further. There is a shift in the phase of the oscillations as a function of $p$, which is consistent with a small shift of frequency, equivalent to the Fermi surface area changing by ~10% between 0 and 0.84 GPa (Fig. 3B). Again, the change is largest between 0 and 0.37 GPa and then varies linearly with further increase in pressure. On depressurization to $p = 0$ after the highest pressure, $\rho_c(T)$ for $T > T_c(H = 0)$ approximately returns to the values before pressurization (fig. S3) as does $T_c$, but the magneto-resistance at low temperatures (Fig. 2) remains larger. The phase shift in the depressurized ($p = 0$) state follows the linear relation found at higher pressures (Fig. 3B).

The temperature dependence of the oscillation amplitudes at the different pressures, along with the fits to determine the effective masses, is shown in Fig. 4A. The mass found at low pressures measured inside the pressure cell is the same as the value measured outside the cell where the sample is in contact with liquid helium at low temperatures. This shows that there are no systematic errors in the mass determination introduced by the pressure cell. Figure 4B shows a central result of this paper, namely, that the effective mass strongly decreases with pressure and hence has an inverse correlation with $T_c$.

DISCUSSION

In Fig. 5, we make a comparison between the effective mass and $T_c$ in Y123 (where $T_c$ is varied by changing the oxygen content $\delta$) and in Y124 (where $T_c$ is varied by pressure). The inverse of the mass is plotted to emphasize the almost linear relation between $\rho$ and $n_p$ found in Y123 close to the two putative QCPs (15). It is clear that in Y124, the opposite trend is observed such that the effective mass decreases as $T_c$ increases toward the maximum.

In interpreting these results, it is necessary to discuss how the physical properties of Y124 are changed under pressure and the mechanism of mass enhancement. The main effect of hydrostatic pressure on the structure of Y124 is to decrease all three lattice parameters, but the distance from the CuO$_2$ plane to the apical oxygen reduces twice as rapidly as the c-axis parameter itself (30). This distance is correlated with $n_p$ in Y123 (17, 31) and hence suggests that $n_p$ increases with increasing pressure in Y124 (30). This is supported by experiments showing that, with increasing pressure, room temperature resistivity and thermoelectric power (32, 33) both decrease as does the pseudogap temperature, as measured by NMR Knight shift (34). It is also supported by our observation that the QO frequency in Y124 increases with pressure. In Y123, an approximately linear increase in $F$ has been found with increasing $n_p$ (35), and it is interesting to note that, when plotted against $T_c$, the data for $F(p)$ in Y124 form an almost-perfect linear continuation of the increase in $F$ with $n_p$ for Y123 (fig. S4).

Quantitative estimates of $dn_p/dp$ (Supplementary Materials) and the fact that the maximum $T_c$ in Y124 and Y123 under pressure is ~105 K (which is considerably higher than the maximum $T_c$ of ~94 K that can be achieved by changing $\delta$ in Y123 alone) suggest that charge transfer is not the only pressure-induced effect that increases $T_c$. It has recently been proposed that an important additional effect could be weakening of CDW order with pressure (36). Because the CDW competes for states with superconductivity, as evidenced by the observed dip in $T_c$ and upper critical field $H_{c2}$ at the value of $n_p$ where the CDW is strongest (37), $T_c$ may be enhanced if the CDW is weakened. Although weakening of the CDW order remains a possibility at higher pressures, we find no decrease in the amplitude of the QO with increasing $p$ (Fig. 3A). This strongly suggests that there is no significant decoherence of the high-field CDW for $p < 0.84$ GPa. The increase in QO amplitude with increased $p$ indicates an increase in the CDW coherence.

As $T_c$ increases linearly with $p$ ($<4$ GPa), it is unlikely that the anomalously large effect of the initial pressurization on $F$ (Fig. 3B)
and the magnitude of the magnetoresistance (Fig. 2) are mainly due to charge transfer. More likely, it results from an increase in the order of any residual vacancies in the material. This would lead to an increase in the mean free path and, in turn, the QO amplitude and the size of the magnetoresistance. When the pressure is removed, this increased order may persist, giving rise to the observed hysteresis of $F(p)$ (Fig. 3B) and the magnetoresistance. However, we find that this increased order does not have any significant effect on $m^*$ (Fig. 4) and so does not affect our central conclusion.

The increase in $m^*$ in Y123 close to optimal doping (15) could be interpreted as resulting from enhanced coupling to fluctuations in either CDW or pseudogap order because these modes would be expected to soften close to their end points. Our key finding is that in Y124, $m^*$ decreases as $T_c$ increases under pressure, which clearly demonstrates that the enhancement found for $m^*$ in Y123 as $\delta$ is reduced, increasing $T_c$ toward optimal, is not universal. Rather, it suggests that the approximate coincidence of $n_{p,c2}$ (the doping where $1/m^*$ extrapolates to 0 in Y123) and the doping where $T_c$ is maximal are accidental, so there is no causal connection between the fluctuations giving rise to the mass enhancement and $T_c$. The decrease in $m^*$ in Y124 could be interpreted as resulting from a decrease in CDW fluctuations caused by the pressure-induced structural changes. For example, the reduction in $c$-axis length caused by increasing pressure might help stabilize the high-field CDW phase where three-dimensional coupling is important and hence increases the CDW coherence, thereby reducing the mass enhancing fluctuations. Elucidation of the microscopic changes will require a structural study of the CDW at high pressure. Our results suggest that the proximity of the CDW end point to the maximum in $T_c$ with doping is coincidental and that, therefore, quantum fluctuations of the CDW order do not boost $T_c$ in cuprates.

**MATERIALS AND METHODS**

**Crystal growth and characterization**

Crystals of Y124 were grown from a mixture of Y123 powder and CuO$_2$ in a KOH flux (38). The obtained crystals had a typical size of $200 \times 200 \times 100$ $\mu$m, with the shortest dimension being along the $c$ axis. The composition was checked using x-ray diffraction, and the lattice parameters were found to be in agreement with previous reports. The samples displayed a sharp superconducting transition with $T_c = 79$ K. For resistance measurements, four contacts were made (two each on the top and bottom faces) using evaporated gold pads, which were annealed onto the sample at 500°C and then contacted with DuPont 4929 silver epoxy. Typical contact resistances were <0.1 $\Omega$ at 80 K.

**Fig. 4. Quasiparticle effective mass in Y124 under hydrostatic pressure.** (A) Amplitude of QOs as a function of temperature. The curves have been offset vertically for clarity. The solid lines are fits to the Lifshitz-Kosevich formula. The field windows used are given in table S1. (B) Variation of quasiparticle mass with pressure extracted from the fits. The dashed line is a guide to the eye. a.u., arbitrary units.

**Fig. 5. Quasiparticle mass of Y124 compared to Y123 plotted versus $T_c$.** The top scale shows the hole doping level for Y123 (17). Data for Y123 are taken from Ramshaw et al. (15) and the references therein. The dashed line is a guide to the eye.
4.2 K. In zero field, resistance was measured using a standard low-frequency lock-in amplifier method, and in pulsed field, resistance was measured using a numerical low-digitization technique typically operating at 54 kHz. Resistance-versus-temperature curves are shown in fig. S3B. Rₐ shows metallic behavior at low temperatures and is similar to the results reported by Hussey et al. (39).

**Pressure cell**

A nonmagnetic pressure cell was used in the pulsed magnetic field system (29). Because of space limitations inside this cell, pressure was determined using the samples themselves as a manometer. Changes in the Rₐ values of the Y124 samples were measured in the nonmetallic cell at room temperature, and these results were compared to measurements performed in a different larger cell where the pressure was measured at low temperatures using a Pb manometer and at room temperature using the resistance of a manganin wire (40). In this second cell, it was verified, by c-axis resistivity measurements, that dTᵣ/dP was 5.5(5) K/GPa, in agreement with previous measurements (20).

**Effective mass determination**

QO data were analyzed by fitting the following function to the Rₓ(B) data

\[ R(B) = A B^{0.5} \exp(-\alpha/B) \sin(2\pi F/B + \phi) + a_0 + a_1 B + a_2 B^2 + a_3 B^3 + a_4 B^4 \]  

The first part of the function represents the field-dependent part of the Lifshitz-Kosevich formula (41) for the QO amplitude in a three-dimensional metal at T = 0. We fit this to a short section of data (field ranges for each pressure are given in table S1) at the lowest temperature (typically T = 1.4 K) and determine α, F, and ϕ. These parameters are then left fixed, and the oscillation amplitude A was then determined as a function of temperature. A fourth-order polynomial (a₀ to a₄) was used to subtract the nonoscillatory background resistance. The temperature dependence of A was then fitted to the thermal damping factor Rₓ = X / sinh(X), X = 2π²kₖm*²T/(ehB) to extract the effective mass m* in this formula, B is set to the average inverse field in the field range fitted.

**Thermometry in the pulsed magnetic field**

To accurately determine the mass, it is necessary to accurately measure the sample temperature. This is not trivial in a pulsed-field system because the relatively high sample currents needed to obtain a sufficient signal-to-noise ratio produce heating at the sample contacts. These effects were minimized in our measurements by obtaining very low contact resistances (around 0.1 Ω) and having the sample in direct contact with liquid helium (for T ≤ 4.2 K) for measurements outside of the pressure cell.

For measurement inside the pressure cell, the thermometer was mounted directly outside the cell. Tests using a miniature GeAu thermometer mounted inside the cell (attached to the sample) showed that the two thermometers were in thermal equilibrium provided that the resistive heating at the sample contacts was sufficiently low. The maximum current used for the resistivity measurements was selected so that this was always the case.

The strong increase in the H of the superconducting-normal transition with temperature was used as a second check of the temperature. The irreversibility field Hᵣₑₑ, defined as the field where the resistance first increases above the noise, varies linearly with T (fig. S5). The continuity of this linear behavior when the sample is inside the cell provides an important check that the sample temperature has been correctly determined.

**Estimates of change in doping with pressure**

A quantitative estimate of the change in n_p with p can be made using the magnitude of the room temperature a-axis thermoelectric power (S_a(300 K)). Using the data for Y124 in the study by Zhou et al. (32) along with the universal relation between n_p and S_a(290 K) (33, 42), we obtained dnp/dp = 2.3(6) × 10⁻³ holes/GPa. A second estimate can be made using the pressure-dependent NMR Knight shift (31) K for Y124 from Meissner et al. (34). Following Tallon et al. (43), we scaled the temperature axis by the pseudogap energy E_g so that the data for different pressures collapse onto a single curve and thus is approximately constant for T/E_g > 1. For p = 0, we estimated E_g = 350 K, which is consistent with the a-axis resistivity ρ_a(T) data (39). From this, we obtained dnp/dp = 3.6(4) × 10⁻³ holes/GPa for p ≤ 4.2 GPa, which is in reasonable agreement with the above estimate from S_a(290 K).

Taking the average of these two values and combining them with an estimate of dE_g/dp, taken from Y123 at n_p = 0.13 (17) would predict that for Y124, dE_g/dp = 2.7(5) K/GPa [that is, about two times smaller than the measured dE_g/dp = 5.5(1) K/GPa] (23).

Although the quantitative precision of this estimate should be viewed with caution because the universal relations between S_a(290 K) and n_p (33, 42) and between E_g and n_p (43) may not hold under finite pressure, it is likely that there are other effects that increase T_c in addition to charge transfer.

In simple systems, the QO frequency F, which is proportional to the Fermi surface extremal cross section, can often be used to precisely determine the doping level. However, in cuprates, this has so far only been successfully demonstrated in the overdoped cuprate Tl₂Ba₂CuO₆+δ (44–46). For underdoped materials, the reconstruction of the Fermi surface complicates the analysis because the area of the reconstructed pocket will depend on the unknown details of the reconstruction potential, as well as n_p. In Y123, an approximately linear increase in F with increasing n_p (35) has been found. If we assume that dF/dn_p in Y124 is the same as in Y123 for 0.09 < n_p < 0.125, then the slope dF/dp from the linear part of Fig. 3B would give dnp/dp = 1.0(2) × 10⁻³ holes/GPa, which is around three times larger than the estimate from S_a(300 K). However, because it remains unknown how universal E_F(nₚ) is in cuprates, this estimate is probably the most uncertain.

**SUPPLEMENTARY MATERIALS**

Supplementary material for this article is available at http://advances.sciencemag.org/cgi/content/full/2/3/e1501657/DC1

Fig. S1. Sample dependence of zero-pressure QOs.

Fig. S2. Effective mass determination at p = 0 for two further samples.

Fig. S3. Pressure dependence of ρ_a(T).

Fig. S4. Comparison of changes in QO frequency between Y123 and Y124.

Fig. S5. Temperature and pressure dependence of the irreversible field H_irr.

Table S1. Field windows used for the fitting of temperature-dependent QO amplitudes at various pressures, as shown in Fig. 4.

**REFERENCES AND NOTES**

1. S. Sachdev, B. Keimer, Quantum criticality. Phys. Today 64, 29–35 (2011).

2. L. Taillefer, Scattering and pairing in cuprate superconductors. Annu. Rev. Condens. Matter Phys. 1, 51–70 (2010).
3. T. Shibaiuchi, A. Carrington, Y. Matsuda, A quantum critical point lying beneath the superconducting dome in iron pnictides. Annu. Rev. Condens. Matter Phys. 5, 113–135 (2014).

4. D. M. Broun, What lies beneath the dome? Nat. Phys. 4, 170–172 (2008).

5. M. Hücke, N. B. Christensen, A. T. Holmes, E. Blackburn, E. M. Forgan, R. Liang, D. A. Bonn, W. N. Hardy, G. Utokowski, M. v. Zimmermann, S. M. Hayden, J. Chang, Competing charge, spin, and superconducting orders in underdoped YBa2Cu3O6. Phys. Rev. B 90, 054514 (2014).

6. S. Blanco-Canosa, A. Frano, E. Schierie, J. Porras, T. Loew, M. Minola, M. Bluschke, E. Weschke, B. Keimer, M. Le Tacon, Resonant x-ray scattering study of charge-density wave correlations in YBa2Cu3O6. Phys. Rev. B 90, 054513 (2014).

7. J. L. Tallon, J. W. Loram, The doping dependence of Tc — What is the real high-Tc phase diagram? Physica C 349, 53–68 (2001).

8. A. Damascelli, Z. Hussain, Z.-X. Shen, Angle-resolved photoemission studies of the cuprate superconductors. Rev. Mod. Phys. 75, 473–541 (2003).

9. T. Wu, H. Mayaffre, S. Krämer, W. N. Hardy, R. Liang, D. A. Bonn, M. H. Julien, Magnetic-field-induced charge-stripe order in the high-temperature superconductor YBa2Cu3O6. Nature 477, 191–194 (2011).

10. G. Ghiringhelli, M. Le Tacon, M. Minola, S. Blanco-Canosa, C. Mazzoli, N. B. Brookes, G. M. De Luca, A. Frano, D. G. Hawthorn, F. He, T. Loew, M. M. Sala, D. C. Peets, M. Salluzzo, E. Schierie, R. Sutarto, G. A. Sawatzky, E. Weschke, B. Keimer, L. Braicovich, Long-range incommensurate charge fluctuations in (Y,Nd)Ba2Cu3O6. Science 337, 821–825 (2012).

11. J. Chang, E. Blackburn, A. T. Holmes, N. B. Christensen, J. Larsen, J. Mesot, R. Liang, D. A. Bonn, W. N. Hardy, A. Waterpufin, M. v. Zimmermann, E. M. Forgan, S. M. Hayden, Direct observation of competition between superconductivity and charge density wave order in YBa2Cu4O8. Nat. Phys. 8, 871–876 (2012).

12. T. Wu, H. Mayaffre, S. Krämer, M. Horvatic, C. Berthier, W. N. Hardy, R. Liang, D. A. Bonn, M.-H. Julien, Incipient charge order observed by NMR in the normal state of YBa2Cu3O7. Nat. Commun. 6, 6438 (2015).

13. S. Gerber, H. Jang, H. Nojiri, H. Matsuzawa, D. A. Bonn, R. Liang, W. N. Hardy, Z. Islam, A. Mehta, S. Song, M. Sikorski, D. Stefanescu, Y. Feng, S. A. Kivelson, T. P. Devereaux, Z.-X. Shen, C.-C. Kao, W.-S. Lee, D. Zhu, J.-S. Lee, Three-dimensional charge density wave order in YBa2Cu3O7 at high magnetic fields. Science 350, 949–952 (2015).

14. M. Nardone, A. Audouard, D. Vignolles, L. Brossard, Hydrostatic pressure anvil cell for electrical transport measurements to 1 TPa under high pressure. Phys. Rev. B Condens. Matter 87, 047007 (2013).

15. B. Vignolle, D. LeBoeuf, D. Vignolles, M.-H. Julien, From quantum oscillations to charge order in high-Tc copper oxides in high magnetic fields. C. R. Phys. 14, 39–52 (2013).

16. M. G. Delucchi, A. Carpentieri, R. A. Cooper, M. M. French, A. P. Mackenzie, C. Jaudet, D. Vignolles, E. Schierle, B. Keimer, M. Salluzzo, E. Weschke, R. Sutarto, G. A. Sawatzky, E. Weschke, B. Keimer, M. Le Tacon, Resonant x-ray scattering study of charge-density wave structural changes of superconducting YBa2Cu4O8 under high pressure. Physica C 450, 023705 (2014).

17. J. D. Jorgensen, B. V. Veal, A. P. Paulikas, L. J. Nowicki, G. W. Crabtree, H. Claus, W. K. Kwok, Structural properties of oxygen-deficient YBa2Cu3Oy superconductors. Phys. Rev. B 84, 1863–1877 (1990).

18. C. Castellani, C. Di Castro, M. Grilli, Non-Fermi-liquid behavior and d-wave superconducting state in YBa2Cu3O7 near optimal doping. Phys. Rev. B 56, 3372–3375 (1997).

19. D. A. Bonn, W. N. Hardy, A. Watenphul, M. V. Zimmermann, E. M. Forgan, S. M. Hayden, D. A. Bonn, W. N. Hardy, R. Liang, S. Adachi, N. E. Hussey, B. Vignolle, C. Prout, M. Sutherland, S. Krämer, J.-H. Park, D. Graf, N. Doiron-Leyraud, L. Taillefer, Direct measurement of the upper critical field in cuprate superconductors. Nature 539, 51 (2016).

20. G. L. Sun, Y. T. Song, C. T. Lin, Investigation of YBa2Cu3O6 single crystal growth by KOH flux. Supercond. Sci. Technol. 11, 66 (2000).

21. N. E. Hussey, K. Nozawa, H. Takagi, S. Adachi, K. Tanabe, Anisotropic resistivity of YBa2Cu3O7. Incoherent-to-metallic crossover in the out-of-plane transport. Phys. Rev. B 56, R11423(9) (1997).

22. C.-y. Wang, Electrical resistance of magnesium coil to 7 kbar and 200°C. Rev. Sci. Instrum. 38, 24 (1967).

23. D. Schoenborn, Magnetic Oscillations in Metals (Cambridge Univ. Press, Cambridge, 1984), 570 pp.

24. J. L. Tallon, C. Bernhard, H. Shaked, R. L. Hitterman, J. D. Jorgensen, Generic superconducting phase behavior in high-Tc cuprates: Tc variation with hole concentration in YBa2Cu3O6. Phys. Rev. B 63, 220517 (2001).

25. J. L. Tallon, J. W. Loram, G. V. M. Williams, J. R. Cooper, I. R. Fisher, D. J. Johnson, M. P. Staines, C. Bernhard, Critical doping in overdoped high-Tc superconductors: A quantum critical point? Phys. Status Solidi B 215, 531–540 (1999).

26. B. Vignolle, A. Carrington, R. A. Cooper, M. M. French, A. P. Mackenzie, C. Jaudet, D. Vignolles, C. Prout, N. E. Hussey, Quantum oscillations in an overdoped high-Tc superconductor. Nature 455, 952–955 (2008).

27. A. F. Bangura, P. M. C. Rouke, T. M. Bensen, M. Matsuishi, J. R. Cooper, N. E. Hussey, A. Carrington, Fermi surface and electronic homogeneity of the overdoped cuprate superconductor Tl2Ba2CuO6 as revealed by quantum oscillations. Phys. Rev. B 82, 140501 (2010).

28. P. M. C. Rouke, A. F. Bangura, T. M. Bensen, M. Matsuishi, J. R. Cooper, A. Carrington, N. E. Hussey, A detailed de Haas–van Alphen effect study of the overdoped cuprate Tl2Ba2CuO6. New J. Phys. 12, 105009 (2010).

Acknowledgments: We thank J. Cooper, S. Friedemann, J. Tallon, and L. Taillefer for useful discussions. Funding: This work was supported by the Engineering and Physical Sciences Research Council (grant no. EP/K016709/1). A portion of this work was performed at the Laboratory National de Champs Magnétiques Intenses, which is supported by the French National Research Agency (ANR) SUPERFIELD, the European Magnetic Field Laboratory, and the LABEX NEXT. We also acknowledge the support of the High Field Magnet Laboratory, Radboud University Fundamental Research on Matter. Author contributions: The experiment
was conceived by A.C., N.E.H., C. Putzke, and C. Proust. QO experiments were performed in Toulouse by C. Putzke, L.M., S.B., D.V., W.T., A.C., and C. Proust. The samples were grown and characterized by P.W., M.B., and C. Putzke. The paper was written by A.C. and C. Putzke with input from all of the co-authors. **Competing interests:** The authors declare that they have no competing interests. **Data and materials availability:** All data needed to evaluate the conclusions in the paper are present in the paper and/or the Supplementary Materials. Additional data related to this paper may be requested from the authors. Raw resistance versus field or resistance versus temperature data for all figures can be found at doi:10.5523/bnrs.p6wv9g3eoh3t1fd2d4fujuno.

Submitted 17 November 2015
Accepted 28 January 2016
Published 18 March 2016
10.1126/sciadv.1501657

**Citation:** C. Putzke, L. Malone, S. Badoux, B. Vignolle, D. Vignolles, W. Tabis, P. Walmsley, M. Bird, N. E. Hussey, C. Proust, A. Carrington, Inverse correlation between quasiparticle mass and $T_c$ in a cuprate high-$T_c$ superconductor. *Sci. Adv.* 2, e1501657 (2016).
Inverse correlation between quasiparticle mass and $T_c$ in a cuprate high-$T_c$ superconductor

Carsten Putzke, Liam Malone, Sven Badoux, Baptiste Vignolle, David Vignolles, Wojciech Tabis, Philip Walmsley, Matthew Bird, Nigel E. Hussey, Cyril Proust and Antony Carrington (March 18, 2016)

Sci Adv 2016, 2.
doi: 10.1126/sciadv.1501657

This article is published under a Creative Commons license. The specific license under which this article is published is noted on the first page.

For articles published under CC BY licenses, you may freely distribute, adapt, or reuse the article, including for commercial purposes, provided you give proper attribution.

For articles published under CC BY-NC licenses, you may distribute, adapt, or reuse the article for non-commercial purposes. Commercial use requires prior permission from the American Association for the Advancement of Science (AAAS). You may request permission by clicking here.

The following resources related to this article are available online at http://advances.sciencemag.org. (This information is current as of February 20, 2017):

Updated information and services, including high-resolution figures, can be found in the online version of this article at:
http://advances.sciencemag.org/content/2/3/e1501657.full

Supporting Online Material can be found at:
http://advances.sciencemag.org/content/suppl/2016/03/15/2.3.e1501657.DC1

This article cites 44 articles, 4 of which you can access for free at:
http://advances.sciencemag.org/content/2/3/e1501657#BIBL

Science Advances (ISSN 2375-2548) publishes new articles weekly. The journal is published by the American Association for the Advancement of Science (AAAS), 1200 New York Avenue NW, Washington, DC 20005. Copyright is held by the Authors unless stated otherwise. AAAS is the exclusive licensee. The title Science Advances is a registered trademark of AAAS