Robustness of the far-field response of nonlocal plasmonic ensembles

Contrary to classical predictions, the optical response of few-nm plasmonic particles depends on particle size due to effects such as nonlocality and electron spill-out. Ensembles of such nanoparticles are therefore expected to exhibit a nonclassical inhomogeneous spectral broadening due to size distribution. For a normal distribution of free-electron nanoparticles, and within the simple nonlocal hydrodynamic Drude model, both the nonlocal blueshift and the plasmon linewidth are shown to be considerably affected by ensemble averaging. Size-variance effects tend however to conceal nonlocality to a lesser extent when the homogeneous size-dependent broadening of individual nanoparticles is taken into account, either through a local size-dependent damping model or through the Generalized Nonlocal Optical Response theory. The role of ensemble averaging is further explored in realistic distributions of isolated or weakly-interacting noble-metal nanoparticles, as encountered in experiments, while an analytical expression to evaluate the importance of inhomogeneous broadening through measurable quantities is developed. Our findings are independent of the specific nonclassical theory used, thus providing important insight into a large range of experiments on nanoscale and quantum plasmonics.

Plasmonics lies among the most prominent research fields in modern nanotechnology, promising exciting applications and unravelling new phenomena as the length scale reduces. Traditionally, noble metals constitute the material basis for novel plasmonic devices operating in the visible, although many recent efforts are devoted to extensions towards the ultraviolet, infrared and THz parts of the spectrum. A key issue in noble-metal plasmonics is its association with pronounced homogeneous broadening due to Ohmic losses in the metal and enhanced Landau damping near the surface. Within classical electrodynamics, and in the quasistatic regime, radiation losses are small and the limited quality factor of plasmon resonances reflects material losses. In other words, homogeneous broadening is important. Furthermore, the commonly employed local-response approximation (LRA) of classical electrodynamics predicts size-independent resonances for the nowadays experimentally accessible small nanoparticles (NPs) in the quasistatic regime. As a consequence, despite the increasing impact of plasmonics and the promotion of single-particle spectroscopy, little, if any, emphasis has been placed on the role of inhomogeneous broadening due to size distribution — even in experiments on NP ensembles with a noticeable size variation.

The observation of size-dependent resonance shifts not anticipated from classical electrodynamics has recently renewed interest in plasmons in the sub-10-nm regime. State-of-the-art experiments range from single-particle spectroscopy with the aid of tightly focused electron beams, to optical far-field measurements sampling the response of NP ensembles. In the latter case, nonlocal effects and the concomitant inhomogeneous broadening can prove important for the interpretation of ensemble measurements. Ensemble averaging effects have been theoretically explored for exciton systems, and for large-NP plasmonic collections dominated by retardation-driven redshifts, but related studies in nonlocal plasmonics are still missing. The unambiguous observation of size-dependent resonance shifts in single-particle spectroscopy encourages therefore to explore broadening phenomena related to size distribution: What is the robustness of plasmonic nonlocal effects when subject to ensemble averaging?

The influence of ensemble spectral averaging on the far-field response of nonlocal plasmonic NP collections is studied here theoretically, starting with the ideal case of a normal distribution of free-electron, Drude-like nanospheres. Complexity is subsequently increased by considering more realistic distributions, resembling experimental histograms, of noble-metal NPs, for which additional loss mechanisms like interband transitions and
Results and Discussion

Nonlocality-induced plasmon blueshifts. We first revisit the optical response of a small metallic nanosphere, embedded in air for simplicity. Our study is based on Mie theory and its appropriate extension for nonlocal effects (see the Methods section for more details). The metal is described as a free-electron plasma with transverse ($\varepsilon_t$) and longitudinal ($\varepsilon_l$) dielectric function components given by the frequency- ($\omega$) and wave vector- ($\mathbf{k}$) dependent Drude and hydrodynamic models, respectively

$$
\varepsilon_t(\omega) = \varepsilon_{\infty} - \frac{\omega_p^2}{\omega(\omega + i\gamma)}, \quad \varepsilon_l(\omega, \mathbf{k}) = \varepsilon_{\infty} - \frac{\omega_p^2}{\omega(\omega + i\gamma) - \beta^2k^2},
$$

where $\omega_p$ is the plasma frequency of the metal, $\varepsilon_{\infty}$ is the background contribution of bound electrons and ions, $\gamma$ is the damping rate, $\beta$ the hydrodynamic parameter, and $k_l$ the longitudinal wave number. We take $\varepsilon_{\infty} = 1$ and $\gamma = 0.01 \omega$, to focus on the role of free electrons and ensure low loss (associated with homogeneous broadening). We further assume $\beta = \sqrt{3/5}v_F$, as obtained within the Thomas–Fermi theory, where $v_F$ (taken equal to $1.39 \times 10^6$ m s$^{-1}$ in the rest of the paper) is the Fermi velocity of the metal.

The size dependence of the frequency of the first (dipolar) plasmonic mode sustained by such a metallic nanosphere of radius $R$ is plotted in Fig. 1a as obtained within LRA ($\omega_{LRA}$, red line) and the hydrodynamic Drude model (HDM) ($\omega_{HDM}$, blue line). To make our results scalable for different materials, $\omega$ and $R$ are normalised to the plasma frequency and wavelength, $\omega_p$ and $\lambda_p = 2\pi/c\omega_p$, respectively. For a better illustration of the sizes and energies usually encountered, the corresponding plasmon energy (NP radius) is provided at the top (right) axis, assuming a typical value $\hbar\omega_p = 9$ eV$^{35}$. For very small NP sizes, LRA reproduces the quasi-static result, $\omega_{LRA} = \omega_p/\sqrt{3}$ (vertical dashed line in Fig. 1a). For larger sizes, retardation causes the modes to drastically redshift and become wider, as also observed in the normalised extinction cross section ($\sigma_{ext}$) spectra of Fig. 1b (red lines corresponding to different NP sizes within LRA). Higher-order modes will not concern us here, and the quadrupolar plasmon peak of the largest sphere in Fig. 1b is only shown by thin dotted lines. The small-size modal frequency saturation predicted by LRA gives place to a continuous blueshift when the metal nonlocal response is taken into account. Comparison between LRA and HDM (blue lines in Fig. 1b) immediately shows that the frequency shifts become larger as the NP size decreases, but no additional resonance broadening due to nonlocality is observed.
needs merely to replace Generalized Nonlocal Optical Response (GNOR) theory. Within SDD, the damping parameter should suffice for the study of ensemble averaging, regardless of the direction and origin of modal shifts.

**Inhomogeneous broadening in Drude-metal ensembles.** Ensemble spectral averaging is at a first step investigated by considering a collection of N = 1000 of the NPs described above, with a mean diameter \( 2(R)/\lambda_p = 0.031 \) (corresponding to 4.3 nm for \( h\omega_p = 9 \) eV). The NP size follows normal distributions around this mean value as shown in the inset of Fig. 2, with standard deviations ranging from 0.2 (narrowest distribution, solid line) to 0.4 (dashed line) and 0.6 (widest distribution, dotted line). The extreme case of a \( \delta \)-function distribution, i.e., all NP diameters corresponding precisely to the mean value, is depicted by open dots. This kind of \( \delta \)-function distribution is exactly what one assumes in practice when disregarding ensemble averaging. We also note that, while the distributions of Fig. 2 are continuous functions, discrete size steps are taken in the simulations, small enough to achieve convergence of the averaged spectra. Apart from the LRA and HDM models, we also discuss calculations based on the commonly employed local size-dependent damping (SDD) model and the Generalized Nonlocal Optical Response (GNOR) theory. Within SDD, the damping parameter \( \gamma \) becomes size dependent, \( \gamma \rightarrow \gamma + 4\pi v_F/\lambda_p \), to effectively take into account the experimentally observed single-NP damping. The constant \( A \), usually taken equal to 1 (as we do here) although a large range of values can be found in literature, is introduced to phenomenologically describe the reduction of the free-electron path length and to account to some extent for quantum-size corrections in very small NPs. On the other hand, GNOR reproduces size-dependent damping in a more physical way, by incorporating electron diffusion as a measure of a variety of electron-scattering effects, including Landau damping due to generation of electron-hole pairs. In practice, one needs merely to replace \( \beta \) in equation (1) with \( \beta^2 + D(\gamma - \omega) \), where \( D \) is the diffusion constant of the metal. A thorough discussion on the determination of \( D \) can be found in a recent review by Raza et al.; in general, it has to be chosen so as to reproduce the experimentally observed, and successfully reproduced by SDD models, plasmon damping. For the Drude-like NPs studied in this section, we find that the simple relation \( D \approx v_F^2/\gamma^2 \) provides an excellent correspondence between the two models. However, it has been shown that more strict calculations are required in the case of noble metals. While, in this respect, GNOR remains a phenomenological model, its strength is that, for arbitrarily shaped plasmonic NPs, it reproduces both the size-dependent blueshifts and the damping of plasmon modes by a simple correction in the dynamics of the free-electron fluid of HDM, whereas SDD models only capture the damping effects.

With these models at hand, we study in Fig. 2 how spectral averaging compares to single-NP response. Clearly, for the local models (LRA and SDD, red and black lines respectively), averaging does not practically affect the spectra. For all size distributions, the average extinction \( \langle \sigma_{ext} \rangle \) normalised to the geometrical cross section of the mean-size NP, \( \pi(R)^2 \) (which is known in experiments), reproduces almost perfectly the spectrum of the individual mean-size NP, without frequency shifts or line broadening. Comparison with Fig. 1 shows that, in the size range of interest, local theories have already reached the quasistatic limit and the plasmon frequencies do not shift

Figure 2. Averaged normalised extinction \( \langle \sigma_{ext} \rangle \) spectra calculated for \( N = 1000 \) NPs described by the dielectric function of equation (1) within the LRA (red lines), HDM (blue lines), GNOR (green lines) and SDD (black lines) models, for the size distributions shown in the inset. The average NP diameter is \( 2(R)/\lambda_p = 0.031 \), which for \( h\omega_p = 9 \) eV corresponds to 4.3 nm, and the standard deviation of the normal distribution function is 0.2 (solid lines), 0.4 (dashed lines), and 0.6 (dotted lines). Open circles denote the corresponding spectra for the single mean-size NP, corresponding to the \( \delta \)-function distribution (open circles) of the inset.
Figure 3. Averaged normalised extinction spectra calculated for $N = 1000$ silver NPs described by the experimental dielectric function of Johnson and Christy\textsuperscript{45} within the LRA (red line), HDM (blue line), GNOR (green line), and SDD (black line) models, for the size distribution shown by the histogram of the inset. The mean NP diameter is $\langle R \rangle = 4.3$ nm. Open circles denote the corresponding spectra for the single mean-size NP.

Inhomogeneous broadening in noble-metal ensembles. The important result of negligible effect of spectral averaging when single-NP size-dependent damping is taken into account may be appealing, but its validity was displayed only for ideal Drude metals and for normal size distributions. In order to connect with more practical, experimentally feasible situations, it is therefore important to carry out similar statistical studies for more realistic distributions in noble metals. We consider a collection of $N = 1000$ silver NPs, described by the experimental dielectric function ($\varepsilon_{\text{exp}}$) of Johnson and Christy\textsuperscript{45}, following the size distribution shown by the histogram of the inset of Fig. 3. In order to apply the HDM, SDD and GNOR models, we obtain $\varepsilon_{\infty}$ in equation (1) from the experimental values by subtracting the Drude part: $\varepsilon_{\infty} = \varepsilon_{\text{exp}} + \omega_p^2/\omega(\omega + \gamma)$, taking $\hbar \omega_p = 8.99$ eV and $\hbar \gamma = 0.025$ eV, values which describe bulk silver excellently. For SDD and GNOR we further assume $A = 1$ and $D = 3.4/10 A v_F^2/(5\gamma_p)$, respectively\textsuperscript{55}. The calculated spectra of Fig. 3 display now an almost negligible difference between single-NP and averaged spectra, even for the more pronounced in Fig. 2 HDM case. Homogeneous line broadening dominates the ensemble optical response, especially when single-NP size-dependent damping is taken into account within the more complete GNOR theory. This observation further strengthens our conclusion that inhomogeneous line broadening is not pronounced in most realistic NP ensembles (despite the non-negligible nonlocal response). Far-field optical experiments on small-NP ensembles can indeed be conducted for the observation of nonlocal frequency shifts, and their interpretation can be performed on the basis of the properties of the mean-size NP in the collection.

In addition to the study of isolated NPs, and the displayed robustness of their far-field optical response, an aspect that we have so far disregarded is the interaction between NPs in the ensemble. It is widely known that once plasmonic NPs are brought close to each other their interaction leads to significant modal redshifts, increasing as the interparticle gap is reduced\textsuperscript{46}. A modified optical response is therefore expected for an ensemble of interacting small NPs, where two competing mechanisms, those of size-dependent blueshifts and interaction-induced redshifts are simultaneously present. The importance of this interplay is explored here, assuming that the average NP distance does not become smaller than $R$, thus preventing the particles from entering the nearly-touching regime, where purely quantum effects such as tunnelling become relevant\textsuperscript{47,48}. Such distance control can be achieved nowadays with unprecedented precision, in dilute solutions with DNA- or ligand-functionalised NPs\textsuperscript{50,59-51}. We assume $N_g = 1000$ dimers of identical, 4.3-nm silver NPs, separated by a gap of width $d$, as shown schematically in the inset of Fig. 4. The interparticle gap width follows a normal distribution around its mean value, $\langle d \rangle = 3.2$ nm, ranging from 4.3 nm (a full NP width separating the two spheres) to 2.1 nm, as shown by the histogram of the inset. Comparison between Figs 3 and 4 shows that NP interaction can lead to a small plasmon redshift, of about 3–4 nm, both for the local case (only LRA is shown in Fig. 4 as the effect of SDD is well reproduced by its nonlocal
Figure 4. Averaged normalised extinction ($\langle \sigma_{ext} \rangle/2\pi R^2$ here, since the geometrical cross section corresponds to the area occupied by two NPs) spectra calculated for $N_p = 1000$ silver NP dimers (both NPs have a diameter of 4.3 nm), separated by a gap of width $d$, as shown schematically in the inset, within the LRA (red line), HDM (blue line), and GNOR (green line) models. The gap width follows the normal distribution of the histogram of the inset, with a mean value $\langle d \rangle = 3.2$ nm. Open circles denote the corresponding spectra for the single mean-gap dimer.

Analytical evaluation of the importance of inhomogeneous broadening. Having considered situations where inhomogeneous broadening can be either strong or negligible, a simple way to decide on its importance without resorting to detailed simulations is desirable. To this end, we develop an analytical model which describes inhomogeneous broadening in terms of just the first two negative-order (or, with some further approximations, positive-order) moments of any NP-size distribution function. In practice, with a simple experimental size histogram at hand, one should be immediately able to tell whether the spectra are affected by inhomogeneous broadening. We begin by considering the dipole resonance in a single metallic NP, neglecting homogeneous broadening for the moment. Such a resonance can then be described by a spectral function $F(\omega, R) \approx \delta(\omega - \omega_{\text{LRA}}) - \eta/R$, where $\eta(\propto \beta$ in our case) gives the strength of the leading-order $1/R$ correction associated with nonlocal response. In an ensemble of non-interacting particles characterised by a size distribution $P(R)$, the ensemble-averaged spectral function will be $\langle F(\omega) \rangle = \int dR F(\omega, R) P(R)$. Our aim is to express the ensemble-averaged optical properties, such as the resonance frequency $\langle \omega \rangle$, with the aid of the $n$th-order statistical moments of the particle ensemble, i.e. $\langle R^n \rangle = \int_0^\infty dR R^n P(R)$. The homogeneous delta-function line shape allows to express the nth-order spectral moment $\langle \omega^n \rangle = \int d\omega \omega^n F(\omega)$ directly in terms of moments of the particle-size distribution,

$$\langle \omega^n \rangle = \int dR (\omega_{\text{LRA}} + \eta/R)^n P(R) = \langle (\omega_{\text{LRA}} + \eta/R)^n \rangle.$$

It is then straightforward to derive expressions for $\langle \omega \rangle$ and the inhomogeneous broadening width, $\Delta \omega_{\text{inhom}} = \sqrt{\langle \omega^2 \rangle - \langle \omega \rangle^2}$, through the statistical moments of the particle-size distribution. As a key result, which allows to estimate the inhomogeneous broadening only in terms of the first two statistical moments of $P(R)$ and the nonlocal blueshift $\omega_{\text{LRA}} - \omega_{\text{NL}} = \langle \omega \rangle - \omega_{\text{LRA}} = \eta(\propto \beta) \langle R^{-1} \rangle$ (and $\langle R \rangle$ in a more crude approximation), it is shown that

$$\frac{\Delta \omega_{\text{inhom}}}{\omega_{\text{LRA}} - \omega_{\text{NL}}} = \sqrt{\frac{\langle R^{-1} \rangle^2 - \langle R \rangle^2}{\langle R^{-1} \rangle^2}} \approx \sqrt{\frac{\langle R^{-1} \rangle}{\langle R \rangle}}.$$

The first equality relates to the first and second negative-order moments of $P(R)$, which are quite unusual ways of characterising a particle-size distribution – in most other contexts the positive-order moments (such as the mean value and variance) are the ones of interest. In the Discussion of the Supplementary Information we demonstrate the link between negative- and positive-order moments to obtain the second approximate identity in equation (3), which links directly to the relative particle-size fluctuation $\Delta R/\langle R \rangle$. This result holds for any description beyond classical electrodynamics that gives a $1/R$ leading-order blueshift of the LRA resonance.
Increasing point size schematically depicts increasing distribution width. The black line denotes FWHM of the distribution function in an ensemble. Our work provides therefore a valuable, general tool for the analysis of far-field optical spectra in modern experiments on plasmonics. While ensemble averaging is shown to produce almost negligible deviations in most situations of practical interest, as illustrated for realistic size distributions of noble-metal NPs, and within the more accurate GNOR model. Nanoscale experiments involving large numbers of NPs can thus be designed and analysed in terms of the averaged mean-size NP in the ensemble, while far-field spectra of large NP collections are still expected to display the fingerprints of nonlocality, as in single-particle spectroscopies. We derived a simple equation to directly identify whether inhomogeneous broadening becomes important, simply through knowledge of the size distribution of the ensemble. The extinction cross section of these NPs, which scales with the distribution width increases from 0.13 · 10^{-2} to 2.80 · 10^{-2} (0.18 nm to 3.86 nm), while the standard deviation of the normal distribution (red circles) increases from 0.32 · 10^{-2} to 7.00 · 10^{-2} (0.04 nm to 0.965 nm). Increasing point size schematically depicts increasing distribution width. The black line denotes FWHM = Δω_{inhom} + FWHM_{0}.

Figure 5. Parametric plot (open symbols) of Δω_{inhom} calculated from equation (3) versus FWHM obtained from simulations for a Drude-like NP within HDM (β = \sqrt{3/5} v_F, γ = 0.01 ω_p in equation (1), for the size distributions shown in the inset. The average NP diameter is fixed at 4.3 nm when hω_{p} = 9 eV). Three different size distributions are plotted: uniform (blue line), triangular (green line) and (truncated) normal (red line). For the uniform (blue squares) and triangular (green triangles) cases, the distribution width increases from 0.13 · 10^{-2} to 2.80 · 10^{-2} (0.18 nm to 3.86 nm), while the standard deviation of the normal distribution (red circles) increases from 0.32 · 10^{-2} to 7.00 · 10^{-2} (0.04 nm to 0.965 nm). Increasing point size schematically depicts increasing distribution width. The black line denotes FWHM = Δω_{inhom} + FWHM_{0}.

frequency. Most importantly, it does not change if we replace η with −η to describe a corresponding 1/R redshift, so that our findings can be easily generalised to include other nonclassical effects, as anticipated above.

To test the validity of equation (3), we use it to evaluate Δω_{inhom} for certain distribution shapes and widths, assuming for simplicity η = β. The result is then compared to the full-width-half-maximum (FWHM) of the (averaged) plasmon peak calculated in each case by simulations performed for an ideal free-electron metal within HDM, with β = \sqrt{3/5} v_F and γ = 0.01 ω_p. As long as equation (3) holds, the average NP size considered in Fig. 5 corresponds to 4.3 nm, while the standard deviation of the normal distribution (red circles) increases from 0.32 · 10^{-2} to 7.00 · 10^{-2} (0.04 nm to 0.965 nm). Increasing point size schematically depicts increasing distribution width. The black line denotes FWHM = Δω_{inhom} + FWHM_{0}.

Conclusion
In summary, the effect of inhomogeneous broadening of plasmon resonances due to nonlocal response in ensembles of small plasmonic NPs was explored through detailed simulations and analytical modelling. While inhomogeneous broadening is negligible in the LRA, it can be an important issue for Drude-like metals, especially within the standard HDM approach which neglects size-dependent damping in individual NPs. Crucially, ensemble averaging is shown to produce almost negligible deviations in most situations of practical interest, as illustrated for realistic size distributions of noble-metal NPs, and within the more accurate GNOR model. Nanoscale experiments involving large numbers of NPs can thus be designed and analysed in terms of the response of the mean-size NP in the ensemble, while far-field spectra of large NP collections are still expected to display the fingerprints of nonlocality, as in single-particle spectroscopies. We derived a simple equation to directly identify whether inhomogeneous broadening becomes important, simply through knowledge of the size distribution function in an ensemble. Our work provides therefore a valuable, general tool for the analysis of far-field optical spectra in modern experiments on plasmonics.
Methods
Nonlocal Mie theory. Here we summarise the fully-retarded Mie theory for a spherical plasmonic particle treated within HDM (and GNOR through a simple substitution of the hydrodynamic β parameter). The multipolar response of a sphere including nonlocal effects was determined by Ruppin[29,33] by extending Mie theory to take into account the excitation of longitudinal waves. In the framework of Mie theory, the extinction cross section of a sphere of radius \( R \) embedded in a homogeneous host medium is given by[13]

\[
\sigma_{\text{ext}} = -\frac{2\pi}{k_h^2} \sum_{\ell=1}^{\infty} (2\ell + 1) \text{Re}(t_{\ell}^{\text{TE}} + t_{\ell}^{\text{TM}}),
\]

where \( \ell \) denotes the angular momentum and \( k_h \) is the wave number in the host medium, which is described by a dielectric function \( \varepsilon_h \). Assuming that the magnetic permeabilities, both in the sphere and in the host medium are equal to 1, the nonlocal Mie scattering coefficients are[29,33–35]

\[
t_{\ell}^{\text{TE}} = -\frac{j_\ell(x_1)}{j_{\ell+1}(x_1)} \left[ x_1 j_{\ell+1}(x_1) \right]' + \frac{j_{\ell+1}(x_1)}{j_\ell(x_1)} \left[ x_1 j_\ell(x_1) \right]' - \frac{\varepsilon_{\text{in}}}{\varepsilon_{\infty}} \frac{j_\ell(x_1)}{j_{\ell+1}(x_1)} \left[ x_1 j_{\ell+1}(x_1) \right]' + \Delta_\ell,
\]

\[
t_{\ell}^{\text{TM}} = -\frac{\varepsilon_{\text{in}}}{\varepsilon_{\infty}} \frac{j_\ell(x_1)}{j_{\ell+1}(x_1)} \left[ x_1 j_{\ell+1}(x_1) \right]' - \frac{\varepsilon_{\text{in}}}{\varepsilon_{\infty}} \frac{j_{\ell+1}(x_1)}{j_\ell(x_1)} \left[ x_1 j_\ell(x_1) \right]' + \Delta_\ell,
\]

where \( j_\ell(x) \) and \( h_\ell^+(x) \) are the spherical Bessel and first-type Hankel functions, respectively, while \( x_1 = k_h R \) and \( x_2 = k_t R \). Here \( k_t \) is the (transverse) wave number inside a sphere described by a transverse dielectric function \( \varepsilon_t \).

The nonlocal correction \( \Delta_\ell \) to the Mie coefficients is given as

\[
\Delta_\ell = \ell(\ell + 1) f_{\ell} (x_1) \frac{\varepsilon_{\text{in}} - \varepsilon_{\infty}}{\varepsilon_{\infty}} \frac{j_\ell(x_1)}{j_{\ell+1}(x_1)},
\]

where \( x_1 = k_t R \) and \( k_t \) is the longitudinal wave number in the sphere, associated with the longitudinal dielectric function \( \varepsilon_t \), which is frequency- and wave vector-dependent. The dispersion of longitudinal waves is given by \( \varepsilon_\ell(\omega, k) = 0 \). In the limiting case where \( \Delta_\ell = 0 \) we retrieve the local result of standard Mie theory. All numerical results for isolated NPs have been obtained from numerical evaluations of equation (4). The corresponding results for NP dimers were obtained by use of a commercial finite-element method solver (COMSOL Multiphysics 5.0, RF module), using the appropriate extension to include nonlocal effects.[35,36]

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Author Contributions
All authors contributed to the conception of the project and writing of the manuscript. C.T. and J.R.M. performed the simulations, C.T., J.R.M., M.W. and N.A.M. analysed the results.

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