Repulsion of polarised particles from anisotropic materials with a near-zero permittivity component

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Reduction of adhesion and stiction is crucial for robust operation on nanomechanical and optofluidic devices as well as atom and molecule behaviour near surfaces. It can be achieved using electric charging, magnetic materials or light pressure and optical trapping. Here we show that a particle scattering or emitting in close proximity to an anisotropic substrate can experience a repulsive force if one of the diagonal components of the permittivity tensor is close to zero. We derive an analytic condition for the existence of such repulsive force depending on the optical properties of the substrate. We also demonstrate the effect using realistic anisotropic metamaterial implementations of a substrate. The anisotropic metamaterial approach using metal–dielectric and graphene–dielectric multilayers provides a tuneable spectral range and a very broad bandwidth of electromagnetic repulsion forces, in contrast to isotropic substrates.

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INTRODUCTION

The concept of levitation of objects through action at a distance1 and the repulsion of particles from surfaces has great importance in nanotechnology, not only to reduce the unwanted adhesion and stiction of nanomechanical particles2,3 but also to achieve new functionalities in optofluidics and low friction devices4. Beyond trivial examples of repulsion of electric charges and magnetic poles or optical trapping in a focused laser beam, a specific permittivity and/or permeability of the material environment can be used to design electromagnetically facilitated forces. A magnet levitates above a superconductor which has a permeability μ < 0 and, thus, expels magnetic fields from its interior5,6. It has recently been shown that materials with a permittivity (ε) close to zero can be used for levitating an electric dipole7,8, since one of the defining characteristics of such ε-near-zero materials is the absence of electric displacement fields inside the material9. In the latter case, under the quasistatic approximation, valid when the dipole is very close to the surface, an electric dipole experiences an electromagnetic repulsive force away from a homogeneous isotropic substrate with complex relative permittivity ε, when the condition |ε| = |ε' + iε''| < 1 is satisfied7. Such an isotropic ε-near-zero material can be implemented with natural materials near their (fixed) plasma frequencies9,10, but having |ε| ≈ 0 on demand at optical frequencies is challenging due to a limited range of available plasmonic materials and their losses at the plasma frequency11–13.

It has recently been shown that permittivity and permeability not commonly available in nature can be designed using metamaterials. The ε-near-zero properties discussed above have been achieved with both plasmonic and dielectric metamaterials and applied in a wide variety of applications such as phase front tailoring14–16, light supercoupling17–22, optical nanocircuits23, lensing24, cloaking25, non-linearity enhancement26 and many others. Nevertheless, ε-near-zero properties have only recently been realised in the visible spectral range, requiring the use of very elaborate nanostructured three-dimensional metamaterials27,28 limited to micron-sized area samples and fixed light polarisations. At the same time, widely tuneable ε-near-zero behaviour is routinely achieved with anisotropic metamaterials with an effective permittivity tensor given by ε_{	ext{eff}} = \text{diag}(\varepsilon_x,\varepsilon_y,\varepsilon_z) which can be fabricated in macroscopic-size implementations using nanowire composites29 or multilayered slabs30. Some components of the effective permittivity become close to zero at frequencies where the optical properties transition from conventional anisotropic to the so-called hyperbolic dispersion regime, when the real parts of ε_x = ε_y and ε_z have different signs26,31–39.

In this paper, we show that a particle with an induced electric dipole in near-field proximity to a uniaxial anisotropic substrate experiences a repulsive force if one of the components of permittivity is near-zero. We analytically derive the condition for repulsion and confirm the results with the exact electrodynamic calculations for specific examples of stacked slabs implementation of a metamaterial with metal or graphene multilayers. By dropping the requirement of an isotropic substrate, a much less stringent condition is sufficient to achieve repulsion of electromagnetic sources from an anisotropic material. This flexibility makes the proposed levitation suitable for experimental implementation with anisotropic ε-near-zero metamaterials40–47 at microwave and optical frequencies where both nanorod and multilayer composites of macroscopic sizes up to several centimetres are available48,49. The anisotropic metamaterials allow achieving a very broad spectral range where repulsive forces are induced, impossible with isotropic materials.
MATERIALS AND METHODS

We consider the electromagnetic force acting on an electric point dipole source with dipole moment \( \mathbf{p} \) embedded in a medium with isotropic relative permittivity \( \varepsilon_i \) and located at a distance \( h \) above an arbitrary substrate. From the point of view of the dipole, oscillating at a given angular frequency \( \omega \), the knowledge of the substrate is entirely contained in the complex Fresnel reflection coefficients \( r^p(k_i) \) and \( r^s(k_i) \) for p- and s-polarisations, where \( k_i = k_0 \cdot n_i \) and \( k_0 \) is the transverse wave vector of the different spatial components of the dipole’s fields.

The vertical force acting on the dipole can be deduced from the Lorentz electromagnetic force acting on the oscillating charges of a dipole: the time-averaged force \(^{50,51} \) is

\[
\langle F_z \rangle = \frac{1}{2} \left\{ \frac{1}{8 \pi \varepsilon_0 \epsilon_i} \int_{-\infty}^{\infty} k_i \left[ \left( |p_x|^2 + |p_y|^2 \right) (k_i^2 r_i^p - k_i^2 r_i^s) + |p_z|^2 (2 k_i^2 r_i^p) \right] e^{i k_i z_0} d k_i \right\}
\]

where \( k_i = k_0 n_i \) is the wave vector in the upper medium, \( n_i \) is the medium refractive index, and \( k_0 = (k_i^2 - k_s^2)^{1/2} \). The only assumption made was the rotational symmetry of the reflection coefficients \( r^p(k_i) = r^s(k_i) \). The integral in Equation (1) effectively accounts for the force arising from the different spatial components of the reflected fields, spanning both the propagating \( k_i \in [0, k_1] \) and the evanescent \( k_i \in [k_1, \infty] \) components. Equation (1) is an exact expression which is straightforward to numerically integrate. All the figures throughout the paper have been simulated using this exact formulation, but the initial discussion will focus on the quasistatic approximation, allowing analytical treatment in order to get an intuitive understanding of the phenomenon.

The quasistatic approximation can be used when the dipole is very close to the surface of the uniaxial material \( h < \lambda \). In this case, the integral in Equation (1) is dominated by the integrand at very high values of \( k_i = k_0 \), where the reflection coefficients approach constant limiting quasistatic values \( r_i^p \) and \( r_i^s \); the latter is zero in non-magnetic media since \( r_i^s = (\mu_2 - \mu_1) / (\mu_2 + \mu_1) \). This approximation makes the integral analytically solvable, leading to a precise expression for the force at small heights over a non-magnetic medium:

\[
\langle F_z \rangle = -\frac{3}{128 \pi \varepsilon_0 \epsilon_i} \left( |p_x|^2 + |p_y|^2 + 2 |p_z|^2 \right) \text{Re} \left\{ r_i^p \right\} h^4
\]

**Figure 1** Schematic diagram of a dipole above a uniaxial substrate. Electric dipole above a uniaxial anisotropic medium modelled with (a) effective medium theory and (b) stacked slabs implementation.

This force is typically attractive \( (\langle F_z \rangle < 0) \) towards the surface. However, a repulsive force \( (\langle F_z \rangle > 0) \) can be achieved if the condition \( \text{Re} \left\{ r_i^p \right\} < 0 \) can be realised.

RESULTS AND DISCUSSION

**Dipole above uniaxial metamaterial**

In the case of a uniaxial anisotropic substrate with the anisotropic permittivity \( \varepsilon_{\text{eff}} = \text{diag}(\varepsilon_x, \varepsilon_y, \varepsilon_z) \) (Figure 1a), the p-polarised reflection coefficient is given by

\[
r^p(k_i) = \frac{\varepsilon_i k_i}{\varepsilon_k + \varepsilon_i k_i + \varepsilon_k k_i} \rightarrow r_i^p = \frac{\varepsilon_i - \varepsilon_k}{\varepsilon_i + \varepsilon_k} \left( \varepsilon_i / |\varepsilon_i| \right)^{1/2}
\]

where \( k_i = (k_i^2 - k_s^2)^{1/2} \) is the z-component of the wave vector in the upper space, while in the uniaxial substrate it is \( k_{2p} = (k_0^2 \varepsilon_x - (\varepsilon_i / |\varepsilon_i|) k_0^2)^{1/2} \) for the extraordinary wave in the effective medium approximation. Assuming a non-absorbing superstrate (real \( \varepsilon_i \)) and requiring the condition \( \text{Re} \left\{ r_i^p \right\} < 0 \), Equation (3) can be solved analytically for the repulsion of the dipole from the substrate:

\[
|\varepsilon_i| / |\varepsilon_i| < \varepsilon_i^2
\]

This condition ensures the existence of the repulsive force \( \langle F_z \rangle > 0 \) in the quasistatic approximation and provides a simple recipe for designing repulsive force for anisotropic substrates. It should be noted that if a dipole is placed in a medium with a high index material (e.g., water), repulsion is easier to achieve by using a lower-index isotropic dielectric as a substrate (Equation (4)). The real challenge, however, is to achieve repulsion in air or vacuum requiring \( |\varepsilon_i| / |\varepsilon_i| < 1 \), which is the focus of this work. This is a much less strict requirement than for an isotropic substrate requiring \( |\varepsilon_i| < 1 \). In particular, if one of the components of the permittivity tensor is close to zero, \( |\varepsilon_i| \approx 0 \) or \( |\varepsilon_z| \approx 0 \), the value of their respective counterpart components \( |\varepsilon_z| \) or \( |\varepsilon_z| \) can be high without affecting the condition for repulsion.

The modulus operator in Equation (4) works on the complex relative permittivities \( \varepsilon_i = \varepsilon_i + i \varepsilon_i' \) and \( \varepsilon_2 = \varepsilon_2 + i \varepsilon_2' \), so their imaginary part plays an important role. Nevertheless, since the condition given by Equation (4) is not a resonant phenomenon, realistic losses can be easily tolerated. Figure 2a maps the values of \( \varepsilon_i \) and \( \varepsilon_2 \) for which repulsion is achieved for a dipole placed in air. To understand the role of losses, the repulsion condition is plotted for varying imaginary part (for simplicity, both permittivities are assumed to have equal imaginary parts \( \varepsilon_i' = \varepsilon_2' \)). In the case of low losses, there exists a broad range of anisotropic permittivities where the repulsion force exists. With the increase of losses, the range of parameters becomes narrower around the point \( (\varepsilon_i', \varepsilon_2') = (0, 0) \), so that for \( \varepsilon_i' = \varepsilon_2' \leq 1 \) the condition \( |\varepsilon_i| / |\varepsilon_i| < 1 \) is no longer satisfied in the quasistatic approximation for any real values of \( \varepsilon_i \) and \( \varepsilon_i' \).

The exact simulations using Equation (1) show that the repulsive force can be achieved in a broader range of parameters than estimated by the quasistatic condition (Figure 2b and 2c). However, for the increasing height \( h \) of the emitting dipole above the substrate, the force is reduced, following an \( h^4 \) dependence close to the substrate. Away from the substrate, the quasistatic limit (Equation (4)) does no longer exactly apply, and the condition for repulsion turns out to be relaxed, with slightly wider range of required permittivities. Also, notice in Figure 2b that a strongly enhanced force is observed near the surface mode resonance condition \( (\varepsilon_i' / \varepsilon_2' < 0 \) when losses tend to zero), but this enhancement disappears completely with a slight increase in

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Thus, for a given value of filling fraction $f_a$, the condition $|e_{ii}|<e_1^2$ can be mapped directly onto an equivalent condition for the materials’ permittivities of the multilayer $e_1$ and $e_2$ (Figure 3). This parametric plot allows straightforward design of the metamaterial to achieve repulsion in a given spectral range. As an example, we considered the case of the permittivities of silver ($e_1$) and fused silica ($e_2$)—the experimental values were taken from Refs. 56 and 57. By appropriately tuning the filling fraction of silver, different wavelength ranges can be selected to achieve a repulsive force, demonstrating the feasibility to adjust the metamaterial to work at any desired wavelength. For example, a stacked multilayer of silver and silica can achieve a repulsive force in the wide range between 600 nm and 700 nm by choosing a silver filling fraction of 0.1. For smaller thicknesses of the dielectric (larger $f_a$), the requirements on the permittivity of metal are more stringently related to $e_1=0$, while for thicker dielectric layers (smaller $f_a$), a very broad range of negative metal permittivities is acceptable to support repulsive force. The exact force (Equation (1)) calculated in the case $f_a=0.1$ in the effective medium approximation is shown in Figure 3b, and agrees well with the analytically predicted wavelength range.

In order to compare the effective medium theory predictions to the exact solutions for the multilayered composite, the reflection coefficient of the semi-infinite structure of stacked slabs in Equation (1) needs to be evaluated. The reflection coefficients of a multilayer can be analytically calculated using the well-known transfer-matrix method for 2N slabs (N unit cells) on a substrate55. The reflection coefficient of a semi-infinite stack of two alternating layers can be calculated as an eigenvalue problem55. We confirmed that, in all cases, increasing the number of unit cells converged towards the semi-infinite stack limit. As detailed in Ref. 55, the correspondence between the multilayered material and the equivalent effective medium breaks if the dipole is placed too close to the surface. In this case, the most rapidly decaying near fields dominate the force (large $k_f$). For these $k_f$, the evanescent fields decay so rapidly that they no longer “see” the stacked structure, and instead feel only the first material slab. In this case, the reflection coefficient approaches the quasistatic value of the reflection coefficient of the first interface between the superstrate $e_1$ and the first slab $e_2$. Therefore, it is highly dependent on the material placed first55. Thus, in the quasistatic approximation $k_f \rightarrow \infty$, valid when the dipole is very close to the semi-infinite stack, the dipole will only interact with the first interface (material $e_2$), resulting in an attractive force (because $|e_1|>|e_2|$) as shown in Figure 3c. The value of $k_f$ for which the stack is a good approximation of the effective uniaxial medium can be increased if the unit cell of the multilayer is made smaller. Increasing the distance from the dipole to the surface attenuates the evanescent components with high $k_f>k_f^{\text{max}}$. Thus, in order to achieve a repulsive force, we face a trade-off when using stacked slabs: either the dipole should be sufficiently far from the composite, so that only spatial components with fields that penetrate several layers dominate.
or the unit cell is sufficiently thin to allow the strongly decaying near fields to penetrate several layers. This trade-off is clearly visible in Figure 3c. Nevertheless, in a reasonable range of small dipole heights and realistically big unit cells (Figure 3b), the force is repulsive and converges to that of the effective material model as the unit cell becomes smaller. Interestingly, the force is higher compared to the predictions of the effective medium model. For thicker unit cells, the force decreases but its bandwidth increases significantly (Figure 3b).

### Increasing repulsion bandwidth

Isotropic ε-near-zero materials typically work in narrow bandwidths around specific frequencies. However, by properly engineering anisotropic metamaterials, we can achieve extremely large bandwidths of the anisotropic ε-near-zero property. The trend in Figure 3a for the repulsive force condition |ε_a|/|ε_t| < 1 predicts an ever-increasing bandwidth as f_a approaches zero and ε_t is close to 1. We will focus on achieving |ε_a|≈0 condition, which is sufficient to generate repulsive forces. The in-plane ε_a permittivity component of a stacked structure is simply the average of the constituent materials’ permittivities, weighted by their filling fraction (Equation (5)). Therefore, near-zero ε_a can be easily achieved by combining positive (dielectric) with negative (metallic-like) permittivity materials, as shown above. In order to increase bandwidth, we consider two additional conditions: (i) a negative permittivity material ε_m with a very low filling fraction f_m→0 (e.g., very thin films of a metal or other plasmonic material) and (ii) dielectric layers (filling fraction f_a = 1−f_m) with a very low index ε_d = 1 (nanoporous polymer films of finely controllable thickness and low refractive indices close to 1 can be easily manufactured reproducibly and inexpensively, and are used as antireflection coatings). From Equation (5), the condition |ε_a|<1 can be rewritten as |ε_d f_a + ε_m f_m| < 1. Intuitively, the weighted average results in ε_d being lowered with respect to ε_d≈1, but not too much, thanks to the low filling fraction f_m, resulting in |ε_a|<1. Importantly, the low filling fraction greatly reduces the frequency dispersion in ε_a, achieving a broad frequency range. We can describe this analytically as follows: the condition |ε_d f_a + ε_m f_m| < 1 with ε_d≈1 can be rewritten as 1−2/f_m < ε_m < 1 (this inequality assumes a lossless metal, for simplicity of the argument, but the limiting case is also valid in the general lossy case, as shown later). As one can immediately see, when f_m→0, the condition for |ε_a|<1 becomes simply ε_m(ω)<1, and such permittivity is ubiquitous in plasmonic materials for a wide range of frequencies.

For Drude-like materials, the upper frequency bound is given approximately by the plasma frequency, while the lower frequency bound is limited in practice by the thickness of the negative permittivity material layer and by how close to 1 the permittivity ε_d of the dielectric slabs is. An analysis including losses yields the same conclusion: if we consider that ε_d≥1 is a lossless and dispersionless dielectric and ε_m = ε_m + iε_m, with |ε_m|≥0, then the condition |ε_a|<1 can be written in terms of a required range of values for the real and imaginary parts of ε_m as follows:

\[
|\varepsilon_a| < 1 \Rightarrow \varepsilon_a \left(1 - \frac{1 + \delta}{f_m}\right) \approx \varepsilon_m < \varepsilon_a \left(1 - \frac{1 - \delta}{f_m}\right) \quad \text{and} \quad \varepsilon_m' < \frac{1}{f_m} \tag{6}
\]

where \(\delta = \varepsilon_d^{-1} (1 - (\varepsilon_d f_a)^{-1})^{1/2}\). In the limiting case of low filling factor \(f_a \rightarrow 0\), together with a low index dielectric \(\varepsilon_d = 1\), we obtain the simple condition \(\varepsilon_m < 1\) described above, associated with very high bandwidths for plasmonic materials. This limiting case is true regardless of the losses since \(\lim (1 - \delta) / f_m = 0\) when \(\varepsilon_d = 1\). In the general case, the existence of a lower and upper limit for \(\varepsilon_m\) determines a finite broad spectral range of repulsion, as seen in the following examples.
Figure 4a shows the repulsive force acting on a dipole above a structure of stacked thin gold films (experimental values of permittivity taken from Ref. 56) between nanoporous polymer films (refractive index 1.05) with a very low metal filling fraction of 1/50, for different unit cell thicknesses, and compared to effective medium theory. The bandwidth of the ε-near-zero regime and the associated repulsive force is unusually large, spanning half the visible spectrum and most of the near infrared. In practice, the fabrication of extremely thin metallic films in a stack of nanoporous polymer can be challenging.

As an interesting solution to this problem, we suggest the use of graphene layers, known to be optically equivalent to extremely thin metallic films in a stack of nanoporous polymer can be challenging. The spectra of 2D sheet conductivity of a single sheet of graphene using the Kubo formula, which is known to fit well with experimental results. The scattering rate was set at $\Gamma = 1.29$ meV, which is the highest value from Refs. 59–61, and the results do not depend strongly on the scattering rate. Figure 4c shows the corresponding repulsive force exerted on a dipole above a stack of such graphene layers, sandwiched between layers of 15 nm low-index nanoporous polymer films ($n = 1.05$). We can see a repulsive force in a wavelength range spanning a considerable part of the infrared spectrum (~1.1–2.9 μm), which corresponds to a frequency range 100–270 THz. This constitutes a huge fractional bandwidth well beyond the possibilities of narrowband isotropic ε-near-zero materials, as discussed in detail below.

In the above considerations we modelled two-dimensional graphene using the approach described in the supplementary information of Ref. 59, summarised here. From Maxwell’s equation $\nabla \times \mathbf{H} = \sigma E$, a well-known equivalence between a conductor with volume conductivity $\sigma_v$ and a dielectric with relative permittivity $\varepsilon$ can be derived as $\varepsilon = 1 + i\sigma_v/\omega\varepsilon_0$. A one-atom-thick system such as graphene does not have a well-defined volume conductivity, instead being characterised by a two-dimensional sheet conductivity, but both can be related if we consider a slab of conductor with thickness $t$ in the z-direction. If $t$ is sufficiently small ($t \ll \lambda$) then $E$ is approximately constant with $z$ inside the slab, so the two-dimensional current density traversing a line element across the thickness of the slab is $|J_{zz}| = |\sigma_v E(z)|dz = t \sigma_v E$, equivalent to a sheet conductivity $\sigma_{2D} = \sigma_v t$. This allows us to model graphene as a thin sheet of dielectric with a thickness-dependent permittivity given by $\varepsilon = 1 + i\varepsilon_{2D}/\omega\varepsilon_0$ in the limit of vanishing thickness. It follows that a negative real part in the equivalent permittivity (metallic character) is obtained when the sheet conductivity has a positive imaginary part. In graphene, the frequency up to which this metallic character is achieved is proportional to the chemical potential. As shown in Ref. 59, the thickness dependence in the optical properties cancels out for $t \rightarrow 0$, and the same behaviour as obtained by rigorous calculations using the sheet conductivity model of 2D graphene is recovered. With this model, we can easily calculate the reflection coefficient and forces for stacks of graphene layers and dielectrics in the same way as for metal–dielectric stacks. Using the same analytical transfer matrix method or the eigenvalue problem mentioned for semi-infinite stacks, we can consider any arbitrarily small thickness of graphene. We confirmed that the reflection coefficients and optical forces converged in all our results for the $t \rightarrow 0$ limit. In Figure 4c, we used a graphene thickness of 0.01 nm well inside the convergence region.

**Comparison with isotropic ε-near-zero materials**

Some selected natural materials exhibit isotropic near-zero permittivity at their plasma frequencies or near permittivity resonances. Taking into account the required low loss conditions, repulsion can be observed for silver and silicon carbide at wavelengths of around 320 nm and 10 μm, respectively. This behaviour is fixed at a given wavelength and has a very narrow operating frequency bandwidth. Comparison of the metamaterial structures described above (with anisotropic ε-near-zero behaviour) and natural ε-near-zero materials, silver and hexagonal silicon carbide, clearly shows the increased

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*Figure 4* Repulsive forces with high bandwidth. (a) The spectra of the force acting on a vertical dipole at a height $h = 0.1\lambda$ above an Au/nanoporous polymer ($n = 1.05$) semi-infinite stack metamaterial with $f_{\text{res}} = 0.02 = 1/50$ for different unit cell sizes and for the equivalent homogeneous anisotropic effective medium. The thickness of gold and polymer slabs (in nm) is annotated for each plot. (b) The spectra of 2D sheet conductivity of a single sheet of graphene using the Kubo formula, modelled with a scattering rate $\Gamma = 1.29$ meV and chemical potential $\mu_c = 1$ eV. Regions of positive imaginary part of conductivity result in negative equivalent permittivity. (c) The spectra of the force acting on a dipole at a height $h = 0.1\lambda$ above a graphene/nanoporous polymer ($n = 1.05$) semi-infinite stack metamaterial for $t_{\text{res,cell}} = 15$ nm and the effective medium approximation. The forces are calculated using Equation (1).
In order to account for this effect, we consider a nanoparticle with a polarisability that will also exert a force on it and may influence the repulsion conditions. Scattering from an illuminated small particle, the illuminating light forms a dipole source. Such a model can be applied to a variety of scenarios where the polarised particle is also reflected by the surface, giving rise to standing waves. The dipole-like scattered field \( E_s(r) \) of the polarizable particle is also reflected by the surface, forming a standing wave \( E_{\text{pw}}(r) = E_{\text{pw}} + E_{\text{pw}}' \). The dipole-like scattered field \( E_{\text{pw}}(r) \) is related to the surface dipole moment by the Green function tensor \( G_{ij}(r_0) \) of the scattered-reflected fields of a dipole over a surface. The polarisation vector of the particle is determined by the total incident electric field at \( r \), i.e., \( \mathbf{p} = \chi (E_{\text{pw}}(r) + E_{\text{pw}}'(r)) \), where \( \chi \) is the polarisability of the particle. The matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal matrix inversion is straightforward because \( G_{ij}(r_0) \) is a diagonal

**Figure 5** Comparison of repulsive force from isotropic and anisotropic \( \varepsilon \)-near-zero materials. Spectra of the force per unit radiated power acting on a vertical dipole situated at a height \( h = 0.012 \) above different \( \varepsilon \)-near-zero materials, including natural isotropic materials silver and silicon carbide and three anisotropic metamaterial realisations shown in Figures 3 and 4. Silicon carbide is represented only in labels (the related spectrum at around 10 \( \mu \)m is not shown in order not to obscure other spectra). The fractional bandwidth \( \Delta f / f_0 \) is indicated, where \( \Delta f \) is the repulsion force bandwidth and \( f_0 \) is the centre frequency of the repulsion force spectrum.

**Figure 6** Repulsive forces on illuminated particle. (a) Schematic of a polarisable particle over a surface under plane wave illumination. (b) Different components of the vertical optical force acting on a polarisable particle over an Ag/SiO\(_2\) multilayer metamaterial \( (\varepsilon_0 = 0.1) \) illuminated at a wavelength \( \lambda = 650 \) nm at normal incidence with a power density of 1 W cm\(^{-2}\). Simulations with both the effective medium approximation and the semi-infinite multilayer structure \( (t_{\text{mult.cell}} = 10 \) nm) are presented. The particle polarisability is \( (3.92 \pm 2.37) \times 10^{-32} \) C V\(^{-1}\), corresponding to a core-shell particle with SiO\(_2\) core (inner radius 27.1 nm) and Ag shell (outer radius 30 nm) designed to have a resonant polarisability near 650 nm. The \( (F_z)_{\text{pw}} \) component for the multilayer case is practically identical to that of the effective medium approximation and, therefore, not shown.
Such intensity of illumination can be easily obtained using focusing effective medium theory. In addition, some natural crystals can behave (stacked slabs) not only do not degrade the performance, but actually in fact, we showed that realistic metamaterial implementations fabrication in wide areas up to several centimeters. The repulsion frequencies have already been demonstrated, with straightforward two terms force discussed in the previous sections. Both components of the reflected dipolar scattering. The latter term corresponds to the force due to the illuminating standing wave and the force due to the reflected dipolar scattering. The former, the force is attractive as is typical for the material. However, when the particle is too close to the surface of the stacked slabs, the effective medium approximation fails as the medium considerations and the stacked slabs implementation of the near-field region, as expected, while at larger distances the standing wave can be controlled by appropriately changing the size of the unit cell of the stacked slabs, the effective medium approximation fails as the particle mainly "feels" the presence of the first layer only, and, therefore, the force is attractive as is typical for the material of the top layer. The distance at which the force changes sign in the near-field can be controlled by appropriately changing the size of the unit cell of the metamaterial realisation. The observed repulsive forces are of the order of $10^{-14}$ N for illuminating intensities of 1 kW cm$^{-2}$ (Figure 6b). Such intensity of illumination can be easily obtained using focusing with high numerical aperture objectives [e.g., 25 mW light focused on (50 μm)$^2$].

CONCLUSIONS

We have shown that the condition for levitation of electromagnetic sources near an $\varepsilon$-near-zero material $|\varepsilon_1| \approx 0$ is greatly simplified if we allow the substrate to be anisotropic, by allowing the possibility that only one component of the permittivity tensor is near zero $|\varepsilon_{xx}|, |\varepsilon_{yy}| \approx 0$. This provides freedom for its design in all frequency domains, and readily available anisotropic $\varepsilon$-near-zero materials working at optical frequencies have already been demonstrated, with straightforward fabrication in wide areas up to several centimetres. The repulsion effect is non-resonant and insensitive to material losses. It is also tolerant to the effects of a finite unit cell in a metallo-dielectric stack: in fact, we showed that realistic metamaterial implementations (stacked slabs) not only do not degrade the performance, but actually can improve the repulsion bandwidth and strength compared to the effective medium theory. In addition, some natural crystals can behave as anisotropic $\varepsilon$-near-zero materials near the reststrahlen bands, typically in the mid-infrared spectral region. Although the condition $|\varepsilon_{xx}|, |\varepsilon_{yy}| < \varepsilon_0^2$ is valid under the quasistatic approximation, we showed that taking retardation into account results in an even broader range of permittivity values that provide repulsion. All this amounts to a high robustness of the described repulsion effect. In stark contrast to the characteristic narrowband behaviour of natural isotropic $\varepsilon$-near-zero materials, we proposed a method to achieve an unusually wide bandwidth and tuneable spectral range of repulsion by using thin metallic film stacks or graphene stacks, sandwiched between low-index dielectrics.

Our calculations consider the force on a point electric dipole, thus the results should be applicable to small illuminated particles [scattering light in the Rayleigh limit], radiating fluorescent particles, quantum dots or radiating microwave antennas, where care should be taken to account for their size. Importantly for applications, the described repulsive force on the dipole is a net force (averaged over a period of electromagnetic field oscillations), so the incoherence of radiation from fluorescent particles is not expected to be a problem as long as the coherence time is equal to several oscillations. Very interesting experiments and applications may involve reduced friction expected to occur on particles or atoms scattering near the proposed anisotropic metamaterials. Our calculations are purely classical, but in the optical regime, fluctuation-induced forces (Casimir/van der Waals) should be accounted for, and the wide bandwidth of repulsion presented opens the interesting possibility of a repulsive Casimir force for particles over these structures.

AUTHORS’ CONTRIBUTIONS

FJR-F and AVZ have developed the concept. FJR-F has performed the simulations. Both authors have written the manuscript.

COMPETING INTERESTS

The authors declare no competing financial interests.

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