POSITIONAL COINCIDENCE BETWEEN THE HIGH-LATITUDE STEADY UNIDENTIFIED GAMMA-RAY SOURCES AND POSSIBLY MERGING CLUSTERS OF GALAXIES

Wataru Kawasaki$^{1,2,3}$ and Tomonori Totani$^{4,5}$

Received 2001 August 17; accepted 2002 May 20

ABSTRACT

We report evidence for the first time that merging clusters of galaxies are a promising candidate for the origin of high Galactic latitude, steady, unidentified EGRET gamma-ray sources. Instead of using past optical catalogs of eye-selected clusters, we made a matched filter survey of galaxy clusters over $4^\circ \times 4^\circ$ areas around seven steady unidentified EGRET sources at $|b| > 45^\circ$, together with a 100 deg$^2$ area near the south Galactic pole as a control field. In total, 154 Abell-like cluster candidates and 18 close pairs/groups of these clusters, expected to be possibly merging clusters, were identified within estimated redshift $z_{\text{est}} \leq 0.15$. Five among the seven EGRET sources have one or two cluster pairs/groups (CPGs) within 1$'$ from them. We assess the statistical significance of this result by several methods, and the confidence level of the real excess is maximally 99.8% and 97.8% in a conservative method. In contrast, we found no significant correlation with single clusters. In addition to the spatial correlation, we also found that the richness of CPGs associated with EGRET sources is considerably larger than those of CPGs in the control field. These results imply that a part of the steady unidentified EGRET sources at high latitude are physically associated with close CPGs, not with single clusters. We also discuss possible interpretations of these results. We argue that, if these associations are real, they are difficult to explain by hadronic processes but best explained by the inverse-Compton scattering by high-energy electrons accelerated in shocks of cluster formation, as recently proposed.

Subject headings: cosmology: theory — galaxies: clusters: general — gamma rays: theory — large-scale structure of universe — surveys

1. INTRODUCTION

The Energetic Gamma Ray Experiment Telescope (EGRET) aboard the Compton Gamma Ray Observatory (CGRO) has left us the third EGRET (3EG) catalog (Hartman et al. 1999), the largest and deepest catalog of high-energy gamma-ray sources to date. However, more than 60% of the 3EG gamma-ray sources (170 out of 271) are yet to be identified, mainly because of the poor accuracy of the position determination. The distribution of the unidentified EGRET sources can be explained as the sum of the Galactic ($|b| \leq 40^\circ$) component and another isotropic (likely extragalactic) component (Mukherjee et al. 1995; Özel & Thompson 1996). While several candidates were proposed for the origin of Galactic sources, including molecular clouds, supernova remnants, massive stars, and pulsars (e.g., see Gehrels & Michelson 1999 and references therein), no astronomical object except for active galactic nuclei (AGNs) has been proposed as the origin of the isotropic component consisting of about 20 sources at $|b| > 45^\circ$ (~65 in the whole sky). All AGNs identified as EGRET sources belong to the blazar class, and there is no evidence that other types of AGNs are emitting gamma-rays detectable by EGRET.

Clusters of galaxies have been studied as a possible source of high-energy gamma-rays, since high-energy cosmic rays are expected to exist in intracluster medium (ICM), which could be emitted by member galaxies, or could be generated by AGNs or shocks in cosmological structure formation. Most previous studies concentrated on the hadronic processes, i.e., pion-decay gamma-rays produced by interaction between cosmic-ray protons/hadrons with intracluster matter (Völk, Aharonian, & Breitschwerdt 1996; Berezhinskii, Blasi, & Ptuskin 1997; Colafrancesco & Blasi 1998), and predictions are well below the detection sensitivity of EGRET even for the case of the Coma cluster, for which only an upper limit has been set by EGRET (Sreekumar et al. 1996). However, attention to high-energy emission from nonthermal electrons is recently increasing. Existence of high-energy electrons in intracluster medium has been suggested by diffuse nonthermal hard X-ray emission and diffuse radio emission for several clusters (see, e.g., Sarazin 2002 for a review). Loeb & Waxman (2000) pointed out that the extragalactic gamma-ray background in the GeV band may be explained by the inverse-Compton (IC) scattering of the cosmic microwave background (CMB) photons by electrons accelerated in large-scale shocks generated in structure formation, if about 5% of the shock kinetic energy is converted into nonthermal electrons. Totani & Kitayama (2000, hereafter TK00) calculated the expected gamma-ray source counts by this process and found that a few tens of sources are expected above the EGRET sensitivity from nearby dynamically forming clusters, and a part of unidentified EGRET sources may be accounted for (see also Waxman & Loeb 2000). The preheating of the intergalactic medium, which is inferred from X-ray luminosity versus temperature relation of clusters and groups, may severely suppress the gamma-ray background flux, but still about 10 massive forming clusters could remain as gamma-ray sources detectable by EGRET (Totani & Inoue 2001).

TK00 estimated that “gamma-ray clusters” detectable by EGRET should have a typical redshift of $\lesssim 0.1$ and mass...
of $\sim 10^{15} M_\odot$. However, no statistically significant correlation between the unidentified EGRET sources and known clusters has been found. There are several possible reasons for this result. First, only a small fraction of clusters should be emitting gamma-rays by the process considered by TK00, since the cooling time of electrons emitting high-energy gamma-rays is very short ($\sim 10^6$ yr), and hence only clusters that are dynamically forming with active shocks can emit GeV gamma-rays. The lack of detection from the Coma cluster is thus explained. We may not be able to observe gamma-rays even from clusters with merging signatures in X-ray or radio bands, which remain on much longer timescales than the gamma-ray emission. It is also difficult to select the candidates of gamma-ray clusters from all unidentified EGRET sources, since the Galactic gamma-ray sources extend to relatively high galactic latitude of $|b| \sim 45^\circ$ (Gehrels et al. 2000); a portion of the sources at even higher latitude seem to be variable, and hence they are likely to be AGNs. Even if the candidates are appropriately selected, the typical redshifts reached in the existing all-sky cluster catalogs in optical (Abell, Corwin, & Olowin 1989) or X-ray (Ebeling et al. 1998) are not much greater than the expected redshift of gamma-ray clusters ($z \sim 0.1$), and hence a portion of the gamma-ray clusters detected by EGRET could have been missed by past cluster surveys. These facts make it difficult to search the correlation efficiently. Furthermore, TK00 pointed out that, since gamma-rays can be emitted only from dynamically forming clusters, their structure may be considerably different and extended when compared with stable, well-established clusters detected by X-rays or optical surveys. This effect might make the correlation search with known clusters even more inefficient.

However, most of such forming gamma-ray clusters should have some substructure or merging signature within them, as expected by the hierarchical structure formation in the cold dark matter (CDM) universe. Therefore, an intensive search for these signatures in the regions around the unidentified EGRET sources, with sensitivities better than existing all-sky catalogs of galaxy clusters, is a straightforward test of the gamma-ray cluster hypothesis.

In this paper, we report the first results from our project to systematically examine the gamma-ray cluster hypothesis using optical galaxy data. Among the 19 unidentified EGRET sources at $|b| > 45^\circ$, we focus here on the seven sources classified as “steady” (Gehrels et al. 2000; D. Macomb 2000, private communication) since the remaining 12 variable sources should be other objects such as flaring AGNs. To perform a correlation analysis between the EGRET sources and galaxy clusters more efficiently than past studies, we make a new sample of galaxy clusters detected automatically based on the matched filter cluster finding algorithm (Kawasaki et al. 1998). This catalog should be better for statistical study of correlation than the past optical cluster catalogs selected by eyes that inevitably induce some systematic bias. We found a statistically significant correlation at maximally $3.7 \sigma$ level between the seven EGRET sources and close pairs of galaxy clusters, while no significant correlation was found with single clusters. We will argue that these results give an indirect support, although not conclusive, for the gamma-ray cluster hypothesis.

2. DATA

We use the galaxy sample extracted from the APM catalog. The data around the seven EGRET sources were obtained via APMCAT service, while the data near the south Galactic pole (SGP) were kindly distributed by S. Maddox and M. Irwin. Only blue passband ($O$ or $b_J$) data have been used, since the red passband data seemed much noisier for some EGRET source regions, especially at the edge of photographic plates. Both $O$ and $b_J$ data were available and analyzed for 3EG J1235+0233. After correcting galaxy dimming due to Galactic absorption using the extinction maps and tools by Schlegel, Finkbeiner, & Davis (1998), galaxies within magnitude range of $14 < m_{b_J} < 20$ were selected as the input data for cluster-finder. The seven $4^\circ \times 4^\circ$ areas centered at each EGRET source are searched. For comparison, a 100 deg$^2$ area near the SGP is also used as a control field. Owing to the presence of holes and photographic plate edges in the data region, the total analyzed area is 182.93 deg$^2$.

3. CLUSTER IDENTIFICATION

To make an original cluster sample, we employed a revised version of the matched filter method by Kawasaki et al. (1998), an automated and objective cluster-finding technique based on maximum-likelihood method. Here we briefly describe the essence of this revised matched filter. A likelihood value that a cluster is centered at a given point on the sky is computed using galaxies in a circular region with radius $r_{\text{cr}}$ centered at the point. The circular region is divided into five annular subregions, and the galaxies in each subregion are then used to compare with the “filter,” a model of spatial and magnitude distribution of cluster galaxies. Since we assume the King model for the surface density profile and the Schechter function for galaxy luminosity function, the “filter” has several control parameters including cluster core radius $r_c$, shape parameter of King model $c$, Schechter parameters $M^{*}$ and $\alpha$, redshift $z_{\text{filt}}$, and richness $N'$. $N'$ is defined as the number of cluster galaxies brighter than $m^* + 5$ and within Abell radius ($\sim 1.5 h^{-1}$ Mpc). The relationship between $N'$ and Abell richness $C$ (the number of cluster galaxies within 2 mag from the third brightest galaxy), obtained with Monte Carlo simulation, is given as $C = 1.1 N'^{0.63}$ with uncertainty of 20%. All parameters but $z_{\text{filt}}$, $N'$, and $r_c$ are fixed as $c=2.25$, $\alpha = -1.25$, $M^{*} = -19.44 + 5 \log h$, and $M^{*}_{\text{gal}} = -19.8 + 5 \log h$. For K-correction, the values for E/S0 galaxies by K. Shimasaku (2001, private communication) and Shanks et al. (1984) were used, respectively, for $O$ and $b_J$ passbands. Cosmological parameters are fixed as $h = H_0/100 = 0.8$ and $q_0 = 0.5$. The number and magnitude distribution of the foreground/background galaxies are locally estimated using the galaxies in an annular region around the point with the inner and outer radii of 0.5 and 1$'$, respectively.

---

6 The terms “forming” and “merging” are difficult to clearly discriminate in the standard hierarchical structure formation, and hence we use them in essentially the same meaning.

7 See http://www.ast.cam.ac.uk/~apmcat/.

8 See http://astron.berkeley.edu/davis/dust/index.html.
In the first step of the actual procedure, we fix $(z_{\text{fit}}, r_c)$ as $(0.14, 50 \, h^{-1} \, \text{kpc})$ and tune only $N$ to maximize likelihood at a point in order to simplify calculation. Maximized likelihood and corresponding $N$ are computed at lattice points with the interval of 0.01 to draw a “likelihood map” and a “richness map.” We use the latter to detect clusters because of simpler appearance of clusters in “richness map” (see Kawasaki et al. 1998, Fig. 2). After smoothing the raw “richness map” with Gaussian filter with $\sigma = 0.03$, we detect cluster candidates appearing as local maxima with $N > 161$ (i.e., Abell richness class $\geq 0$). Then $z_{\text{fit}}$ and $r_c$ are surveyed in the range of $0.04 \leq z_{\text{fit}} \leq 0.2$ and $10 \leq r_c \leq 400$ (in $h^{-1} \, \text{kpc}$), respectively, to estimate redshift $z_{\text{est}}$ and richness $N_{\text{est}}$ for each candidate. To avoid erroneous estimation, the above procedure is run for four cases of different galaxy sampling with $r_{\text{cir}} = 0.05, 0.1, 0.15,$ and 0.2. Basically we adopt the values ($z_{\text{est}}, N_{\text{est}}$) for the case $r_{\text{cir}} = 0.2$ unless they are far apart from the other values for the case $r_{\text{cir}} = 0.05, 0.1,$ and 0.15. The uncertainty of $z_{\text{est}}$ is estimated to be $\sim 20\%$ with Monte Carlo simulation. Finally, we obtain a volume-limited, “three-dimensional” sample of 154 cluster candidates with Abell richness class $\geq 0$ complete out to $z = 0.15$.

Figure 1 shows central 2.78 $\times$ 2.68 areas of the “richness maps” around the seven EGRET sources. Cluster candidates are seen as local peaks of color contour. It should be noted that this color contour just indicates the amplitude of the “best-fit” filter with a fixed redshift parameter $(z_{\text{fit}} = 0.14)$ and does not directly reflect cluster’s richness except for the ones at $z = 0.14$. Only the clusters with $z_{\text{est}} \leq 0.15$ and Abell richness class $\geq 0$ (i.e., $N_{\text{est}} > 161$), which we utilize below, are marked with the pluses.

Using this cluster sample, we search for close cluster pairs or groups (CPGs) as candidates of merging clusters. If there are close clusters satisfying the two criteria that (1) their estimated redshifts equal one another within the uncertainty of redshift estimation (20%) and (2) their transverse separation at that redshift is less than 2 $h^{-1}$ Mpc, we regard them as a CPG. In total, we identify 18 CPGs consisting of two to four clusters. Table 1 lists relative position (cols. [3] and [4]) and separation (col. [5]) from the nearest EGRET source, mean estimated redshift (col. [6]), total Abell richness (col. [7]), and number (col. [8]) of member clusters for the nine CPGs found in the EGRET data areas. Some of the CPGs are shown as the green ellipses enclosing member clusters in Figure 1.

4. RESULTS

Here we try several statistical tests for the correlation between clusters and the seven EGRET sources.

4.1. Projected Number Density

We examine if there is an excess overdensity of clusters or CPGs in the vicinity of the EGRET sources (VES). VES is defined as the sum of all areas within 1” of the seven EGRET sources, and the boundary is shown as the yellow solid circles in Figure 1. Considering the extended nature of CPGs, the VES radius is fixed at 1” rather than the EGRET error radius; the value is close to the typical size of both EGRET error circle and expected gamma-ray clusters detectable by EGRET (TK00). The rest of the data area (the EGRET region outside VES plus the SGP region) is hereafter referred to as “control field”. Considering the lack of galaxy data due to the photographic plate edges, VES and the control field cover 20.07 and 162.86 deg$^2$, respectively.

Simply counting all clusters, VES and the control field contain 21 and 133 clusters, respectively. The number of clusters expected by chance in VES should obey the Poisson distribution with the expectation value inferred from the control field, $133 \times 20.07/162.86 = 16.4$ if we ignore cluster-cluster correlation. We see that there is only a weak density excess of clusters at $1.1\sigma$ level in VES.

However, the situation changes greatly for CPGs. Five among the seven EGRET sources, namely, all except for 3EG J1235+0233 and 3EG J1337+5029, have CPGs within 1” from them. Four EGRET sources (3EG J0038–0949, 3EG J1321–1318, 3EG J1310–0517, and 3EG J1347+2932) have one CPG, and the other one (3EG J0159–3603) has two CPGs within 1”. Thus, there are six and 12 CPGs in VES and the control field, respectively. Therefore, the number of CPGs expected by chance in the 20.07 deg$^2$ VES field is $12 \times 20.07/162.86 = 1.5$, thus the number excess of CPGs in VES amounts to $(6 - 1.5)/(1.5)^{1/2} = 3.7\sigma$ level (namely, 99.6% confidence level [CL] assuming Poisson distribution), which is in sharp contrast to the case for single clusters. Even in a conservative case [increasing CPG number of the control field to $12 + (12)^{1/2}$, namely $+1\sigma$ level], the CPG number excess is at $3.0\sigma$ level (or 98.7% CL). The weak correlation between single clusters and EGRET sources seems to appear under the influence of the strong correlation between CPGs and EGRET sources.

4.2. Mean Closest Separation

Next we assess the correlation between CPGs and the EGRET sources in a slightly different way by computing the mean closest separation between CPGs and the EGRET sources and examining if it is smaller than that for the case if CPGs are randomly distributed. Using six EGRET sources except for 3EG J1337+5029, for which no CPG is found in the data area, the mean closest separation is 0.84 arcmin. We then perform a Monte Carlo simulation to compute the mean closest separation for random distribution case. We have 60,000 realizations of random placement of CPGs with density of $12/162.86 = 0.074 / \text{deg}^2$ and then measure the distances of the closest CPGs from a given point. Computing the mean of every six closest separations, we obtain the distribution of 10,000 values of mean closest separation for the random case. The mean and the standard deviation of this distribution are 1.84° and 0.39°, respectively. Using this distribution, the observed mean closest separation for the six EGRET sources is apparently smaller than that for random case with (1.84–0.84)/0.39 = 2.6 $\sigma$ level or 99.8% CL (for the conservative case in the previous subsection, 2.3 $\sigma$ level or 99.5% CL). These results change only very little if we assume that there is a CPG just outside the 4° x 4° area around 3EG J1337+5029.

4.3. Bayesian Statistics Using Elliptical Fits

In addition to the rather simple analyses in §§ 4.1 and 4.2, we also performed more sophisticated correlation study based on the Bayesian statistics, with the same procedure that has been used in some past studies on EGRET source identifications (Mattox et al. 1997; Mattox, Hartman, & Reimer 2001). We can calculate the likelihood ratio (LR) of
Fig. 1.—Matched filter “richness maps” for the seven regions centered at the steady unidentified EGRET sources at $|b| > 45^\circ$. The EGRET source name and the bandpass of the galaxy data are shown at the top of each panel. The plus signs denote the Abell-like cluster candidates with $z_{\text{cl}} \leq 0.15$ and Abell Richness Class $\geq 0$ detected by a matched filter. The small open circles are Abell/ACO clusters for reference. Close cluster pairs or groups (CPGs) are shown as the green ellipses enclosing their member clusters. The VES boundary is shown with large yellow circles (solid line). The yellow dotted ellipses denote the best-fit ellipses for 95% confidence regions of the EGRET sources by Mattox et al. (2001).
identification, \( LR \equiv \frac{dp(r|ID)}{dp(r|c)} \), for a CPG located at a separation angle of \( r \) from the center of an EGRET source, where \( dp(r|ID) \) and \( dp(r|c) \) are differential probabilities that a CPG is found at \( r \) when the CPG is a correct identification of the EGRET source or a confusion noise, respectively. Here, the information of a mean CPG number density \( (12/162.86 \text{ deg}^2) \) and elliptical fits to the 95% CL contour of the likelihood of the EGRET source location (Fig. 1, yellow dotted lines) are used to calculate the LR. (See Mattox et al. 1997 for details.) The distribution of the LR can be used as an empirical indication of the strength of a potential identification. The values of LR are given for the nine CPGs in Table 1.9 To compare with this distribution we performed a Monte Carlo simulation (MC) to produce 900 random locations of nine CPGs (i.e., 100 each) around the six EGRET sources, assuming no correlation between the two. Since the LR of the nine CPGs is distributed in a range of \( 1.4 \times 10^{-6} \) to 9.4, we compare the cumulative distribution of one LR to the MC in the same range, as shown in Figure 2. Clearly, the distribution of the nine CPGs is deviated toward higher LR compared with that of the MC. The Kolmogorov-Smirnov test (KS) gives a chance probability of this deviation as 2.3%; i.e., the observed LR distribution is different from the MC with a confidence level of 97.7%.

Although this result also indicates the physical correlation between the CPGs and EGRET sources, the significance seems less than those estimated in the previous two sections. However, we should emphasize that the test in this section should be conservative for the following reason. Because of the limited time to perform the matched filter calculation, the search for CPGs is made only for regions surrounding EGRET error circles, and hence CPGs far from EGRET sources are excluded in the above sample of CPGs. Therefore, we do not have any real CPG with \( LR < 1.4 \times 10^{-6} \), and we have to compare the observed LR distribution to the MC only in the limited range of \( LR > 1.4 \times 10^{-6} \). This means that the absolute number of CPGs with \( LR > 1.4 \times 10^{-6} \) is not taken into account in the statistical significance. On the other hand, the result of § 4.1 indicates that finding nine CPGs with \( LR > 1.4 \times 10^{-6} \) in the region around EGRET sources is higher than expected from random coincidence. Therefore, the statistical significance only by the KS test in this section might be an underestimate. In addition, we took the separation \( r \) to the center of CPGs, but it is uncertain where the gamma-ray-emitting region is located in the extended region of CPGs. Therefore, the calculation of likelihood ratio might be too strict.

We can infer the a priori probability, \( p(ID) \), that each of the nine CPGs is a correct identification of EGRET sources and the a posteriori probability, \( p(ID|r) \), that a CPG located at \( r \) is the correct identification, as follows. Again, following Mattox et al. (1997), we can calculate \( p(ID|r) \) for each

![Cumulative Distribution](image)

**Fig. 2.**—Cumulative distribution of the likelihood ratio of identification of CPGs as EGRET sources. The solid line is for the nine CPGs around the six EGRET sources considered in this paper, which are listed in Table 1. The dashed line is the result of Monte Carlo simulation assuming no physical correlation between CPGs and EGRET sources. The distribution is considered in a range of \( 1.43 \times 10^{-6} < LR < 9.4 \), which is the range of LR found for the nine CPGs.

---

9 The relative positions of CPGs in Table 1 are from EGRET locations given in the 3EG catalog, while the centers of elliptical fits given by Mattox et al. (2001) are slightly different. We corrected this offset here.

---

**TABLE 1**

| ID (1) | EGRET Source (2) | \( \Delta \alpha \) (deg) | \( \Delta \delta \) (deg) | Separation (deg) | \( z \) (5) | \( C_{\text{total}} \) (6) | \( N_d \) (7) | LR (9) | \( p(ID|r) \) (10) | Note (11) |
|-------|------------------|------|------|----------|-----|---------|------|------|---------|---------|
| 1-1... | 3EG J0038–0949   | 0.85 | 0.26 | 0.88     | 0.055 | 128     | 2    | 6.6E–2 | 2.2E–2  |          |
| 2-1... | 3EG J0159–3603   | 0.04 | −0.64| 0.64     | 0.141 | 109     | 2    | 4.7   | 0.64    |          |
| 2-2... | 3EG J0159–3603   | 0.79 | 0.37 | 0.88     | 0.116 | 206     | 3    | 6.1E–1 | 0.19    |          |
| 2-3... | 3EG J0159–3603   | −1.71| −0.20| 1.72     | 0.104 | 111     | 2    | 1.43E–6| 5.5E–7  | Outside VES |
| 3-1... | 3EG J1234–1318   | −0.35| 0.32 | 0.48     | 0.101 | 217     | 4    | 9.4   | 0.78    |          |
| 4-1... | 3EG J1235+0233   | 0.72 | −1.17| 1.37     | 0.071 | 111     | 2    | 5.7E–3 | 2.2E–3  | Outside VES |
| 5-1... | 3EG J1310–0517   | −0.54| −0.49| 0.73     | 0.092 | 79      | 2    | 2.4   | 0.48    |          |
| 6-1... | 3EG J1347+2932   | 0.91 | −0.21| 0.93     | 0.044 | 165     | 4    | 1.2   | 0.32    |          |
| 7-1... | 3EG J1347+2932   | 0.56 | −1.37| 1.48     | 0.054 | 62      | 2    | 7.5E–2 | 2.7E–2  | Outside VES |
| 8-1... | 3EG J1347+2932   | 0.85 | 0.26 | 0.88     | 0.055 | 128     | 2    | 6.6E–2 | 2.2E–2  |          |

**Note.**—Col. (1): IDs corresponding to those in Fig. 1. Cols. (3)–(5): Relative celestial coordinate and separation from the central position of the 3EG catalog. Col. (6): Mean of estimated redshift for CPGs. Col. (7): Total Abell richness. Col. (8): Number of cluster members included in a CPG. Col. (9): Likelihood ratio of identification as an EGRET source. Col. (10): The a posteriori probability of identification.
observed CPG when an unknown \( p(ID) \) is specified. Then we solve a self-consistent \( p(ID) \) such that the integral of \( p(ID|r) \) divided by the number of CPGs considered (=9) yields the assumed value of \( p(ID) \). We found \( p(ID) = 0.275 \) here, and \( p(ID|r) \) assuming this value of \( p(ID) \) is also given for every CPG in Table 1. According to this \( p(ID|r) \) estimate, we can calculate a chance probability that all CPGs are misidentified, i.e.,

\[
p_c = \prod_{i=1}^{9} \left[ 1 - p_c(ID|r_i) \right],
\]

where the subscript \( i \) runs over the nine CPGs. We found this probability to be 2.2%; i.e., at least one GPG is the correct identification with 97.8% CL.

### 4.4. Estimated Redshift and Richness

The bottom panel of Figure 1 of TK00 shows that the redshift and mass of gamma-ray clusters detectable by EGRET is \( \sim 0.1 \) and \( \sim 10^{15} M_\odot \), respectively. Both of them are roughly consistent with the estimated values of the CPGs in the vicinity of EGRET sources (see Table 1, cols. [6] and [7]). Note that for richness, column (7) of Table 1 shows that the majority of the CPGs have total Abell richness of \( C_{\text{total}} > 100 \) (Abell richness class 2-4). This means that they are quite massive systems with mass of \( \sim 10^{15} M_\odot \).

We also found that the six CPGs within the VES of EGRET sources seem to have larger \( C_{\text{total}} \) compared with those not associated with EGRET sources. The six CPGs have \( C_{\text{total}} = 79, 109, 128, 165, 206, \) and 217, which should be compared with those of 12 CPGs in the control field: 62, 67, 90, 91, 92, 99, 102, 110, 111, 114, 119, and 154. If CPGs are not related to EGRET sources at all, the distribution of richness should be the same for the EGRET region and the control field. The KS test indicates that the chance probability of getting this result from the same distribution function is 8.0%. This is not very compelling by itself, but it should be noted that this test is completely independent of the spatial correlation discussed in the previous three sections. If this result is added to the spatial correlation, significance would be further increased.

### 5. DISCUSSION

#### 5.1. Variability of EGRET Sources

In this work we selected seven sources at \( |b| > 45^\circ \) that are showing no evidence of variability, according to Gehrels et al. (2000). However, the variability of EGRET sources cannot be determined clearly for many cases. In fact, there are two other studies on the variability by Tompkins (1999) and Torres, Pessah, & Romero (2001), and the classification of EGRET sources into variable or nonvariable sources is sometimes different among these authors. We also checked the variability indicators defined by Tompkins (\( \tau \)) and Torres et al. (\( I \)) for the seven EGRET sources here. According to the plausible criteria given by Torres et al. for these two indicators, they can be classified into either “variable,” “dubious,” and “nonvariable” sources. We found that all but one of them are classified as nonvariable or dubious sources in both the two indicators; however, only one source, 3EG J1310–0517 is classified as a variable source by the \( \tau \) indicator, while it belongs to nonvariable sources by the \( I \) indicator. The difference seems to come from the analysis of Tompkins et al. utilizing EGRET data that are not included in the 3EG catalog, while the classification by Gehrels et al. or Torres et al. is based solely on the 3EG catalog (P.L. Nolan 2001, private communication).

Considering this point, we also give statistical significance of correlation when 3EG J1310–0517 is removed from the sample. The number of CPGs expected by chance within the six VES of EGRET sources (18.15 deg\(^2\)) is 1.33, and hence the observed five CPGs are \( (5 - 1.33)/(1.33)^{0.5} = 3.2 \) \( \sigma \) excess of random coincidence (98.8% CL in Poisson distribution). The KS chance probability of the likelihood ratio distribution in the Bayesian analysis becomes 5.1%, and the chance probability that all CPGs are misidentification becomes \( p_c = 7.1\% \).

### 5.2. Theoretical Interpretation and Comments on the Other Work

After the submission of this paper, we learned of a recent study by Colafrancesco (2002, hereafter C02), who investigated the correlation between unidentified EGRET sources and Abell clusters. Our analysis is based on the newly produced cluster catalog based on the automated matched filter method, which is more reliable and objective for statistical cluster study than the eye-selected Abell clusters. On the other hand, C02 also examined radio and X-ray fluxes of clusters associated with unidentified EGRET sources. C02 found interesting correlations between X-ray, radio, and EGRET gamma-ray fluxes that further strengthen the possible connection between clusters and EGRET sources. All but one (3EG J1235+0233) EGRET source considered here are also included in the list of candidates selected by C02. Therefore, at first glance, observational results seem consistent with ours. However, it should be noted that only about half of the CPGs presumably associated to the EGRET sources have Abell clusters as their members. We also found no statistically significant excess of number density for single clusters. The correlation claimed by C02 is between single clusters and a larger number of EGRET sources at \( |b| > 20^\circ \), including more variable sources than considered here. Such a correlation could also be induced by point sources (e.g., AGNs) residing in galaxy clusters. On the other hand, our result that only CPGs show strong correlation with steady EGRET sources indicates that the origin of gamma-ray emission is shocks by cluster/structure formation.

The theoretical interpretation of these results by Colafrancesco is very different from ours; in fact, he strongly argued that the forming/merging clusters proposed by TK00 are not responsible for the association suggested by this work and/or C02. Here we give a detailed interpretation of our results giving some comments on C02’s arguments against forming/merging clusters and argue that the suggestion made by TK00 is the best explanation of the possible association between CPGs and EGRET sources.

To begin with, let us make clear what are the essentially new aspects of the proposal by TK00. This work considers the IC scattering by electrons accelerated in shocks generated by the process of hierarchical cluster formation. This work is the first to predict gamma-ray source counts expected by such process based on the standard structure formation theory in the CDM universe, and TK00 found that maximally a few tens of forming clusters could be detectable by the EGRET. On the other hand, previous
studies concerning gamma-rays from galaxy clusters mostly considered the hadronic processes such as pion decays by primary cosmic-ray protons and emission from secondary electrons (Völk et al. 1996; Berezhiani et al. 1997; Colafrancesco & Blasi 1998). Generally, these papers found gamma-ray fluxes well under the EGRET sensitivity limit, even for a sample of the clusters closest to us, including Coma (Colafrancesco & Blasi 1998). The energy-loss timescale of high-energy protons in clusters is comparable to or longer than the Hubble time, and hence the gamma-ray luminosity should also be steady on this timescale. Then there is no reason to expect even stronger gamma-ray flux from clusters other than Coma. Therefore, gamma-rays produced by cosmic-ray protons in galaxy clusters have not been seriously considered as a candidate of unidentified EGRET sources.

However, high-energy electrons that can emit gamma-rays have a much shorter cooling timescale \( \tau_{\text{cool}} \sim 10^6 \) yr than that of protons or ions. Therefore, if a comparable energy is going into cosmic-ray electrons and protons, then we expect much stronger gamma-ray luminosity of the electron origin when the shock is still active after formation or merging processes, since their energy is emitted within a duration of shock lifetime, i.e., dynamical time \((\sim \text{Gyr})\). (Note that we should not use \( \tau_{\text{cool}} \) here.) Furthermore, we expect gamma-rays only from clusters still having active shocks and do not from well-stabilized clusters without shocks. The gamma-ray flux is expected to vary strongly in the history of hierarchical formation of a cluster. Therefore, it is possible that Coma and other nearby clusters do not have strong gamma-ray emission, while other more distant, or less stabilized clusters emit gamma-rays detectable by EGRET. What TK00 found is that it is, in fact, quantitatively possible, based on the abundance of forming objects in the CDM universe.

C02 first criticized an inconsistency between TK00 and this work: this work sees the correlation between EGRET sources and cluster pairs as the main support of TK00’s idea, although TK00 rested it on forming clusters, which should have more extended galaxy distribution and/or much fainter X-ray flux than normal clusters. It should be noted, however, that only three of the eight CPGs shown in Figure 1 are coincident with Abell clusters. This suggests that a significant part of CPGs found by our paper are very extended and only found as clustering of galaxy clusters, each of which is small and not in the Abell catalog. (However, as shown in §4.4, the total mass of these CPGs is as large as \( 10^{15} M_\odot \).) This is, in fact, consistent with the picture of TK00. On the other hand, coincidence with known Abell clusters for some of CPGs and EGRET sources is also not surprising, since, in the hierarchical structure formation theory, forming or merging clusters should sometimes include rich clusters that can be detected by past surveys. One important point is that, even in such cases, TK00 predicts that gamma-ray emission should not be from the center of rich clusters. Rather, gamma-ray emission is expected from more extended region of CPGs including the rich clusters.

We note that Reimer (1999) set an upper limit on the gamma-ray flux from A85 as \( < 6.77 \times 10^{-8} \) cm\(^{-2}\) s\(^{-1}\) (>100 MeV), instead of accounting for the nearby unidentified source, 3EG J0038−0949. (On the other hand, this association is classified as the most probable association in C02.) If gamma-ray flux is coming from hadronic processes, then we expect that the flux should be the strongest at the center of A85, where cosmic-rays are well confined and ICM density is the highest. Therefore, the hadronic processes cannot explain the association of A85 and 3EG J0038−0949. On the other hand, as discussed by TK00, the IC Compton gamma-ray emission from forming or merging clusters should be more extended because of more extended ICM and uniform density of the CMB. Formation shocks are expected in the surrounding region of the CPG including A85, which seems marginally overlapping with the 95% ellipse of 3EG J0038−0949. Therefore, if A85 and 3EG J0038−0949 are physically associated, the IC Compton gamma-rays should be a better explanation than hadronic processes.

Second, C02 claimed apparent discrepancy between the numbers of clusters in TK00 and this work: 20–50 clusters predicted by TK00 and seven found in this work. Here C02 did not take into account the sky coverage; 20–50 clusters of TK00 are for all sky, but the seven in this work are for \(|b| > 45^\circ\). The correction factor makes these numbers consistent. We also note that 20–50 sources predicted by TK00 might be a rather optimistic value, since electron power index is assumed to be \( \alpha = 2 \) \( (dN_e/dE_e \propto E_e^{-\alpha}) \). Somewhat softer spectrum is expected in reality (see Totani & Inoue 2001), reducing detectable sources by EGRET. However, it should be noted again that even such a reduced number of detectable sources is still much larger than that expected by the hadronic processes considered before TK00.

Third, C02 claimed the difficulties of energetics in the theory of TK00; the gamma-ray luminosity inferred from EGRET sources is much larger than that possible in the TK00’s framework. We again emphasize that TK00 assumes that 5% energy injection from shock kinetic energy to nonthermal electrons. The kinetic energy is calculated based on the standard structure formation theory. It is generally thought that supernova remnants inject about 10% of explosion energy into cosmic-ray protons, although it may seem relatively large for electrons when one considers the energetics ratio of \( \sim 0.01−0.1 \) of cosmic-ray electrons to protons observed at the Earth. However, radio observations of supernova remnants indicate that shocks convert at least a few percent of the shock energy into the acceleration of relativistic electrons (Blandford & Eichler 1987; Sarazin 2002). The extreme-ultraviolet (EUV) and hard X-ray emission from several clusters can be attributed to nonthermal electrons having the total energy of the same order (Fusco-Femiano et al. 1999; Sarazin 2002). Therefore, the energetics assumed by TK00 is not extreme at all. C02 ignored this fundamental point and showed discrepancies in some quantities assuming cluster mass of \( 10^{14} M_\odot \). It is not clear why C02 chose this very small mass instead of the standard value of \( 10^{15} M_\odot \). Since \( L_e \propto M^2 \) in equation (7) of C02, this choice reduces the gamma-ray flux by a factor of 100. On the other hand, TK00’s calculation shows that the typical mass of clusters that are detectable by EGRET is, in fact, \( \sim 10^{15} M_\odot \). The richness estimate of CPGs in this paper also indicates similar masses (see §4.4). Therefore, it is not surprising that C02 found some discrepancy, but it does not give any argument against TK00. C02 also claimed that TK00’s scenario results in an extraordinary temperature of intracluster matter, \( \sim 27−270 \) keV, but it again seems to originate from the nonstandard choice of C02 for cluster parameters \( (T = 8 \text{ keV for } M = 10^{14} M_\odot) \), which is not supported by the \( M-T \) relation of observed clusters [e.g., Finoguenov, Reiprich, & Böhringer 2001] and a small density of
n = 10^{-4} \text{ cm}^{-3} \text{ rather than a typically used value of } 10^{-3} \text{ cm}^{-3}). TK00 predicts that temperature of clusters detectable by EGRET should not be much different from that of normal clusters. Rather, the temperature could be sometimes lower if the shock kinetic energy has not yet been dissipated well in most of the intracluster medium.

Fourth, C02 claimed that there is no evidence for strong ongoing shocks in the sample presented in this work, while the model of TK00 predicts strong nonthermal emission in hard X-ray and EUV bands by the IC scattering, and in the radio band by the synchrotron emission, by the same electron population producing gamma-rays. The TK00 model indeed predicts hard X-ray and EUV emission at a similar flux (in $\nu F_{\nu}$) to that in gamma-ray bands in most cases. However, as repeatedly mentioned by C02, there is almost no observational information in these bands for the sample of EGRET sources considered in this paper. No evidence by no observation is trivial, and it argues against neither TK00 nor this paper. Instead, intensive follow-up observations for the sample presented here in hard X-ray and EUV bands would give an important test of TK00’s scenario. The synchrotron radio flux depends sensitively on the strength of magnetic field; the $\nu F_{\nu}$ flux should scale as $\nu F_{\nu} (\text{radio})/\nu F_{\nu} (\text{gamma}) \sim U_{\text{mag}}/U_{\text{CMB}}$, where $U_{\text{mag}}$ and $U_{\text{CMB}}$ are the energy density of magnetic field and CMB, respectively. When the cluster magnetic field is at a level of $\sim 3 \text{ } \mu \text{G}$ ($U_{\text{mag}} \sim U_{\text{CMB}}$), we expect similar strong flux in radio bands. However, observed hard X-ray flux and radio flux from the Coma cluster indicates $\sim 0.15 \mu \text{G}$ for this cluster (Fusco-Femiano et al. 1999; Sarazin 2002), and then the radio flux should be 0.25% of gamma-ray flux. Considering also that there is no deep radio observations for the sample in this paper, no evidence for strong radio emission does not necessarily contradict with TK00’s suggestion.

C02 noted that A85 in the vicinity of 3EG J0038–0949 is associated with a cold front that is a possible signature of the early stage of merging. C02 claimed that this is not an ongoing violent merging processes, and hence this argues against TK00. However, this is not the case. A85 has a radio halo found on the border between substructures, where the cluster gas is first being shocked (Sarazin 2002). Once the merging starts and shocks are generated, the IC gamma-ray flux should rapidly increase with a minimum timescale of cooling and acceleration of high-energy electrons ($\sim 10^6 \text{ yr;}$ TK00). After this situation is achieved, gamma-ray flux is expected to be rather steady while the shock is propagating. On the other hand, nonthermal radiation in other wave bands (X, EUV, and radio) will achieve this steady state at a timescale of energy dissipation of responsible particles. Since this timescale is longer than that for gamma-ray bands, the initial rise of luminosity could be slower than in gamma-rays. Therefore, strong gamma-ray flux is theoretically expected even when observations in other wave bands show evidences only for early merging stage.

C02 concluded that “the energy release at gamma-ray energies $E > 100 \text{ MeV}$ of the EGRET-cluster associations is probably due to a superposition of diffuse (associated with the active ICM of the cluster) and concentrated (associated with the active galaxies living within the cluster) gamma-ray emission.” We do not disagree with this statement; our claim is that the physical process responsible for “the active ICM of the cluster” should be IC scattering by primary cosmic-ray electrons produced by structure formation, which is the central point of TK00’s proposal. On the other hand, other processes within the standard physics, such as hadronic processes, are unlikely to explain gamma-ray flux detectable by EGRET from active ICM, as mentioned above. In fact, on the basis of the hadronic processes, Colafrancesco & Blasi (1998) predicted gamma-ray flux much smaller than the EGRET sensitivity limits for nearby clusters at $z = 0.01–0.07$. The majority of Abell clusters that are claimed to be associated with EGRET sources by C02 have even larger distances of $z \approx 0.1$ (see C02, Table 1). Since the gamma-ray flux by hadronic processes is not expected to vary significantly from cluster to cluster because of the long dissipation timescale, it seems difficult to explain the gamma-ray flux from ICM by the Colafrancesco & Blasi model.

6. CONCLUSIONS

We performed a correlation analysis between the seven steady unidentified EGRET sources in the high-latitude sky ($|b| > 45^\circ$) and a quasi-three-dimensional catalog of galaxy clusters newly generated with a matched filter algorithm. While there is no correlation between the EGRET sources and the individual clusters, in sharp contrast we found a strong (maximally 99.8% CL level) correlation between the EGRET sources and close pairs/groups (CPGs). This result is consistent with the gamma-ray cluster hypothesis proposed by TK00, which expect that the gamma-ray emission comes from only ongoing mergers with active shocks but not from the usual ones in dynamically “quiet” regimes where the violent shock has subsided. Because of the short timescale of energy dissipation, gamma-ray luminosity should rise more rapidly with the generation of the shock, and decay with the disappearance of the shock, compared with the thermal or nonthermal emission in longer wavelength. This suggests that some clusters may have strong gamma-ray emission with still weak or only early signatures of merging in other bands, and others may have weak gamma-ray flux but with still remaining merging signature in longer wavelength. Confirmation of the merging signatures in CPGs found in this paper is important for further verification, but deep observation is necessary when the merging is still in the early stage.

Clearly, the weak point of our analysis is the small sample (seven) of the steady unidentified gamma-ray sources due to the flux limit of 3EG catalog (although it is the deepest to date). However, TK00 predicted that future gamma-ray telescopes such as GLAST could find hundreds to thousands of gamma-ray clusters up to $z = 0.2–0.3$. The coming three-dimensional deep galaxy catalogs from ongoing SDSS and 2dF survey projects will be ideal resources to directly compare with the GLAST gamma-ray sources. When it is established that a part of the extragalactic steady gamma-ray sources are from forming (merging) clusters, large-scale distribution of gamma-ray clusters will offer unique and valuable information about the dynamical side of cosmological structure formation, in contrast to the more stationary side that has been probed by conventional galaxy clusters in X-ray and optical bands.

---

10 This may not be the case in the very early stage of merging; see the next paragraph.
We are grateful to K. Shimasaku for computing $K$-correction for $O$ passband and $O-\beta y$ color of present E/S0 galaxies. W. K. is supported in part by Japan Society for the Promotion of Science (JSPS) Research Fellowship. T. T. is supported in part by the JSPS Postdoctoral Fellowship for Research Abroad (2001).

REFERENCES

Abell, G. O., Corwin, H. G., & Olowin, R. P. 1989, ApJS, 70, 1
Berezinski, V. S., Blasi, P., & Ptuskin, V. S. 1997, ApJ, 487, 529
Blandford, R. D., & Eichler, D. 1987, Phys. Rep., 154, 1
Colafrancesco, S. 2002, A&A, in press (C02)
Colafrancesco, S., & Blasi, P. 1998, Astropart. Phys., 9, 227
Ebeling, H., et al. 1998, MNRAS, 301, 581
Finoguenov, A., Reiprich, T. H., & Böhringer, H. 2001, A&A, 368, 749
Fusco-Femiano, R., et al. 1999, ApJ, 513, L21
Gehrels, N., Macomb, D. J., Bertsch, D. L., Thompson, D. J., & Hartman, R. C. 2000, Nature, 404, 363
Gehrels, N., & Michelson, P. 1999, Astropart. Phys., 11, 277
Hartman, R. C., et al. 1999, ApJS, 123, 79
Kawasaki, W., Shimasaku, K., Doi, M., & Okamura, S. 1998, A&AS, 130, 567
Loeb, A., & Waxman, E. 2000, Nature, 405, 156
Mattox, J. R., Hartman, R. C., & Reimer, O. 2001, ApJS, 135, 155
Mattox, J. R., Schachter, J., Molnar, L., Hartman, R. C., & Patnaik, A. R. 1997, ApJ, 481, 95
Mukherjee, R., Bertsch, D. L., Dingus, B. L., Kanbach, G., Kniffen, D. A., Sreekumar, P., & Thompson, D. J. 1995, ApJ, 441, L61

Ozel, M. E., & Thompson, D. J. 1996, ApJ, 463, 105
Reimer, O. 1999, in Proc. 26th Int. Cosmic-Ray Conf. (Salt Lake City), 4, 89
Sarazin, C. L. 2002, in Proc. XXI Moriond Conference: Galaxy Clusters and the High Redshift Universe Observed in X-Rays, ed. D. Neumann, F. Durret, & J. Tran Thanh Van (Gif-sur-Yvette: Editions Frontières), in press
Schlegel, D. J., Finkbeiner, D. P., & Davis, M. 1998, ApJ, 500, 525
Shanks, T., Stevenson, P. R. F., Fong, T., & MacGillivray, H. T. 1984, MNRAS, 206, 767
Sreekumar, P., et al. 1996, ApJ, 464, 628
Tompkins, W. 1999, Ph.D. thesis, Stanford Univ.
Torres, D. F., Pessah, M. E., & Romero, G. E. 2001, Astron. Nachr., 322, 223
Totani, T., & Inoue, S. 2001, Astropart. Phys., 17, 79
Totani, T., & Kitayama, T. 2000, ApJ, 545, 572 (TK00)
Volk, H. J., Aharonian, F. A., & Breitschwerdt, D. 1996, Space Sci. Rev., 75, 279
Waxman, E., & Loeb, A. 2000, ApJ, 545, L11