1. INTRODUCTION

The frequency of spectroscopic binaries (SBs) in star-forming regions has been a topic of interest for many years. Mason et al. (1998) review the fraction of binary systems found among O stars in clusters/associations, the field, and runaway systems. The results from their large survey agree with previous studies (e.g., Garmany et al. 1980; Levato et al. 1991) and indicate that a large fraction of O stars in clusters and associations are binaries while fewer field stars show multiplicity and very few runaway stars have companions. The values from Mason et al. (1998) are broken down into spectroscopic and resolved binaries. For O stars in clusters and associations they give a total of 75% as binaries; 61% of the binaries are spectroscopic and 41% are resolved (with some obvious overlap). The spectroscopic binary group includes, however, a significant number of cases labeled as SB1? or SB2? (systems suspected of radial velocity variability or showing double lines), so the fraction of confirmed spectroscopic binaries is then 34% among O stars in clusters and associations.

We have conducted a radial velocity survey of O-type stars in the Cas OB6 stellar association. The Cas OB6 region is one of the largest complexes in the Perseus arm of the Galaxy and is located at a distance of about 2.3 kpc (Garmany & Stencel 1992; Massey et al. 1995; Foster & Routledge 2003). The warm dust and ionized gas in this region extend over 150 pc along the arm and are dominated by the W3/W4/W5 chain of H ii regions and the HB3 supernova remnant (Carpenter et al. 2000; Terebey et al. 2003). The open cluster IC 1805 (OCL 352=Melotte 15) is situated at the western end in the W4 region and contains eight O-type stars (Massey et al. 1995). The three brightest of these, HD 15558, HD 15570, and HD 15629, are probably the ionization source that powers the blowout of gas in the Galactic chimney (Normandeau et al. 1996; Terebey et al. 2003) and the emission arc above W4 (Reynolds et al. 2001). The open cluster IC 1848 (OCL 364) is found in the eastern part of the association in the W5 region. Its primary source of ionizing flux is the multiple O star system HD 17505 (Stickland & Lloyd 2001). In addition to these two clusters, Carpenter et al. (2000) find evidence of 19 other young clusters that are still embedded in their primordial clouds.

Our knowledge of the binary star population among the young, massive stars of Cas OB6 is quite limited at present. There are a number of radial velocity investigations of the sample, but the overall small number of measurements has made the detection of binary stars difficult (Hayford 1932; Underhill 1967; Ishida 1970; Conti et al. 1977; Liu et al. 1989, 1991; Rauw & De Becker 2004). Spectroscopic orbits are available for only three systems, HD 15558 (Garmany & Massey 1981), HD 16429 (McSwain 2003), and BD +60 497 (Rauw & De Becker 2004), while a fourth star, BD +60 470=DN Cas, is a known eclipsing binary (Frazier & Hall 1974). However, results for O stars in other clusters and associations suggest that many more binaries may exist (Mason et al. 1998; Garcia & Mermilliod 2001). It is important to determine the extent of the binary star population in Cas OB6 in order to estimate accurately the ionizing flux of the key O-type stars and to assess the role of gravitational encounters with binary stars that might lead to ejection of runaway stars (Gies & Bolton 1986; Hoogerwerf et al. 2001).
Here we present a reconnaissance of the radial velocity variations observed in 13 bright stars in Cas OB6 that we obtained over a five-night run with the Kitt Peak National Observatory (KPNO) 4 m Mayall telescope in 2003 October. We also present the first radial velocity study of the eclipsing binary system DN Cas, which we combine with published photometry to determine the component masses. We discuss our observations and reduction techniques in § 2. The radial velocity measurements and the derivation of orbital elements are summarized in § 3. We use a tomography algorithm to separate the spectra of the binary star components, and we describe a comparison of the stellar spectra with model spectra in § 4 that we use to estimate the basic stellar parameters. We examine in § 5 the spectral energy distributions to determine reddening, ratio of total to selective extinction, and angular size for each target, and we discuss the probable distance to Cas OB6. Details about the individual objects follow in § 6. Finally, we summarize in § 7 our derived fraction of multiple and single systems for Cas OB6 in the context of past results.

2. SPECTROSCOPIC OBSERVATIONS AND REDUCTIONS

The blue spectra were obtained with the Ritchey-Chrétien (RC) spectrograph at the Mayall 4 m telescope at KPNO from 2003 September 30 through 2003 October 4. The spectra were made with the BL 380 grating (1200 grooves mm⁻¹, blazed at 4500 Å in second order) using a BG39 order sorting filter. The detector was the T2KB CCD, a 2048 x 2048 device with 24 μm square pixels. The spectra recorded wavelengths between 4180 and 4940 Å, although internal vignetting occurs for wavelengths redward of about 4800 Å. The reciprocal dispersion of the spectra is approximately 0.37 Å pixel⁻¹, and the resolving power is R = λ/Δλ = 5700. Exposure times were usually a few minutes in length, and each stellar spectrum was preceded and followed by an He-Ne-Ar lamp exposure for careful wavelength calibration. Most of the resulting spectra have a signal-to-noise ratio (S/N) in excess of 100 pixel⁻¹. We usually obtained at least two spectra of each target on each of five consecutive nights. We also made numerous bias and flat-field frames each night. A list of the names of targets appears in Table 1.

We also obtained an additional 32 spectra of HD 17505A with the KPNO 0.9 m coudé feed telescope in runs in 2000 October and December (Boyajian et al. 2005). The spectra were made with the long collimator, grating B (in second order with order sorting filter OG550), camera 5, and a Ford 3072 x 1024 CCD (F3KB) detector. This arrangement produced a resolving power, R = λ/Δλ = 9530, and covered a range of 829 Å around Hα. Exposure times were 10–15 minutes, and the resulting spectra have an S/N = 200 in the continuum. We also observed the rapidly rotating A-type star, ζ Aql, which we used for removal of atmospheric water vapor lines. Each set of observations was accompanied by numerous bias, flat-field, and Th Ar comparison lamp calibration frames.

The spectra were reduced using standard routines in IRAF to create continuum-rectified versions of each spectrum. Then all the spectra of a given target were transformed to a common heliocentric frame in log λ form and collected in a spectrum matrix of wavelength and time for subsequent analysis. The rectification of the vignetted portion of the spectrum was generally problem-free, so we used the entire spectral range in the following analysis.

3. RADIAL VELOCITIES AND ORBITAL ELEMENTS

We found that the spectra of 10 of the 13 targets were composed of one set of spectral lines, and for these targets we obtained radial velocities by finding the extremum point of the cross-correlation of each spectrum with a spectrum template (omitting those spectral regions containing interstellar features). We initially obtained velocities using the first spectrum as the template, and then we shifted each spectrum to the reference frame of the first spectrum and co-added all the spectra to form a high-S/N template. The final relative velocities were then found by cross-correlation with this high-S/N template spectrum. These relative velocities were then transformed to absolute velocity by

| Object Name       | Spectral Class | \( V_r \) (km s⁻¹) | \( \sigma(u - l) \) (km s⁻¹) | \( \sigma(s - s) \) (km s⁻¹) | \( T_{\text{eff}} \) (K) | log g (cm s⁻¹) | \( V \sin i \) (km s⁻¹) | Binary Status |
|-------------------|---------------|--------------------|----------------------------|----------------------------|-----------------|---------------|----------------------|---------------|
| DN Cas B          | O8 V          | -52                | ...                       | 4.9                        | 149.3           | 34.2          | 4.0                  | 158 ± 20       | SB2            |
| DN Cas B          | B0.2 V        | -38                | ...                       | 5.3                        | 207.3           | 29.7          | 4.0                  | 113 ± 18       | SB2            |
| BD +60 497A       | O6 V          | -43                | ...                       | 2.2                        | 126.8           | 37.5          | 4.0                  | 183 ± 9        | SB2            |
| BD +60 497B       | O8 V          | -38                | ...                       | 3.6                        | 166.7           | 33.0          | 4.0                  | 171 ± 19       | SB2            |
| BD +60 501        | O6.5 V        | -58                | 9.4                       | 5.1                        | 1.1             | 36.8          | 4.0                  | 173 ± 10       | Constant       |
| HD 15558          | O5 III(f)     | -49                | 5.3                       | 1.8                        | 2.9             | 39.5          | 3.7                  | 195 ± 11       | Constant (SB1) |
| HD 15570          | O4 If⁺        | -55                | 0.9                       | 0.8                        | 4.2             | 42.9          | 3.5                  | 123 ± 10       | Atmospheric variations |
| HD 15629          | O5 V(f)       | -60                | 8.1                       | 1.2                        | 2.2             | 41.5          | 4.0                  | 106 ± 6        | Constant       |
| BD +60 513        | O7.5 V        | -59                | 6.7                       | 1.5                        | 2.8             | 34.8          | 4.0                  | 259 ± 5        | Constant       |
| BD +62 424        | O6.5 V(f)     | -34                | 3.7                       | 0.8                        | 1.7             | 36.2          | 3.8                  | 103 ± 12       | Constant       |
| HD 17505Aa        | O7.5 V(f)     | -9                 | ...                       | 8.6                        | 115.3           | 36.3          | 4.0                  | 90 ± 10        | SB2            |
| HD 17505Aa2       | O7.5 V(f)     | -44                | ...                       | 7.2                        | 118.2           | 36.5          | 4.0                  | 120 ± 10       | SB2            |
| HD 17505Ab        | O6.5 III(f)   | -38                | 8.2                       | ...                        | ...             | 37.5          | 3.8                  | 80 ± 15        | Constant       |
| HD 17520          | O9 V          | -54                | 2.7                       | 0.7                        | 6.0             | 33.2          | 4.0                  | 80 ± 12        | SB1 + Be       |
| HD 237019         | O8 V          | -43                | 8.1                       | 0.7                        | 1.1             | 34.0          | 4.0                  | 128 ± 5        | Constant       |
| BD +60 586        | O7.5 V        | -48                | 4.7                       | 1.0                        | 1.0             | 35.4          | 4.0                  | 95 ± 5         | Constant       |
| BD +60 594        | O8.5 V        | -86                | 8.3                       | 1.0                        | 17.7            | 33.5          | 4.0                  | 285 ± 16       | SB1            |

a Systemic velocities for the binary stars are given in Tables 5, 7, and 9.
TABLE 2
RADIAL VELOCITY MEASUREMENTS FOR CONSTANT VELOCITY AND SB1 STARS

| Star Name | Date (HJD - 2,450,000) | $v_r$ (km s$^{-1}$) |
|-----------|-------------------------|---------------------|
| BD +60 501| 2912.891                | -56.6               |
| BD +60 501| 2912.968                | -56.4               |
| BD +60 501| 2913.888                | -57.2               |
| BD +60 501| 2913.965                | -57.9               |
| BD +60 501| 2914.890                | -57.4               |

Notes.—Table 2 is published in its entirety in the electronic edition of the Astrophysical Journal. A portion is shown here for guidance regarding its form and content.

We used a portion of the statistical $F$-test to determine if a star was a velocity-variable object. Since we usually obtained at least two spectra every night of each object and since we do not expect large variations in velocity over the time span of a few hours, we estimated the average measurement error from the mean of the absolute value of the differences within a night divided by $\sqrt{2}$. This spectrum-to-spectrum error estimate $\sigma(s-s)$ is given in column (5) of Table 1. We then compared this to the overall standard deviation of the velocities $\sigma$ (col. [6] of Table 1) and set the criterion for variability as $\sigma/\sigma(s-s) > 3.3$, which corresponds to the <1% probability that both variances are drawn from the same population (for 9 degrees of freedom = 10 observations minus 1 mean value). We found that 3 of the 10 single-lined targets were velocity variable according to this criterion (see col. [10] of Table 1). The stars HD 17520 and BD +60 594 are probably long-period binary systems, while the very luminous star HD 15570 exhibits variations related to atmospheric structure fluctuations (see the discussion in § 6).

The spectra of the targets DN Cas and BD +60 497 both showed evidence of line doubling in some observations, and the spectra of HD 17505A displayed three components in the best-separated cases. We were able to use a tomography algorithm to isolate the stationary line component in HD 17505A ($\xi$ 6.5), and after removal of the stationary component the spectra were treated in the same way as for the other double-lined systems. We measured velocities for these double-lined binaries by interpolating in the O star synthetic spectrum grid of Lanz & Hubeny (2003) for a temperature and gravity appropriate for each star, and these were convolved with broadening functions to account for rotational and instrumental broadening ($\xi$ 4). Then we obtained velocities for specific spectral lines by a nonlinear least-squares fit of the observed profile with the two template spectra profiles co-added with a predefined monochromatic flux ratio (details are given for each case in § 6). We took care to form an average velocity for each spectrum based on the same set of selected lines in order to avoid spurious variations introduced by line-to-line velocity differences. The resulting velocities are given below in § 6 (Tables 4, 6, and 8).

We determined orbital elements for these three double-lined binaries using the nonlinear, least-squares fitting routine of Morbey & Brosterhus (1974). The starting value for the orbital period was obtained with the period search technique of Scargle (1982) as modified by Horne & Baliunas (1986) (or from photometry in the case of DN Cas). We first solved for the best-fit period and then formed separate elliptical solutions for both the primary and secondary components. The derived eccentricity in the case of DN Cas was not significantly different from zero (Lucy & Sweeney 1971), but in the other two targets the eccentricity was significant and the same within errors in the solutions for the primary and secondary components. We formed average values of the eccentricity $e$, longitude of periastron $\omega$, and time of periastron $T$ that were then fixed in solutions to obtain the semi-amplitude $K$ and systemic velocity $\gamma$ for the primary and secondary components individually (since differing atmospheric expansion rates may lead to different systemic velocities). The final orbital elements are given in § 6 (Tables 5, 7, and 9).

4. STELLAR PARAMETERS

We co-added our spectra to form high-S/N versions for use in spectral classification, temperature and gravity estimation, and measurement of the projected rotational velocity $V \sin i$. We used a simple shift-and-add algorithm to form an average spectrum in the case of the single-lined targets. However, we relied on the Doppler tomography algorithm described by Bagnulo et al. (1994) to reconstruct the individual component spectra of the multiple component systems. The tomography algorithm is an iterative scheme that uses the derived orbital velocity shifts and monochromatic flux ratio to estimate the spectra of each component. We ran the algorithm for 50 iterations with a gain of 0.9, although the results are insensitive to both parameters. The final, continuum-rectified spectra of all our targets (except HD 17505; see § 6.5) and all their components are shown in Figures 1 and 2 in order of decreasing effective temperature.

We estimated the spectral classification of each target (Table 1) by comparing their spectra with standard-star spectra in the atlas of Walborn & Fitzpatrick (1990). The corresponding stellar effective temperature $T_{\text{eff}}$ was determined by comparing the equivalent widths of various He I/He II line ratios with those measured in the synthetic spectral library for O stars by Lanz & Hubeny (2003). In most cases, we simply assumed a surface gravity of $\log g = 4.0$ as a suitable value for main-sequence stars, but this approach was inadequate in several cases of evolved stars where the H Balmer line wings are narrower and indicative of a lower gravity. We experimented with a number of model synthetic spectra to best estimate the gravity in these cases. This was accomplished through comparing models by eye with the data in order to find an appropriate match for surface gravity. The most difficult spectrum to match was that of the very luminous star HD 15570, which displays clear evidence of the atmospheric expansion into a stellar wind that is not included in the static atmosphere formulation adopted by Lanz & Hubeny (2003). We estimate that our
fitting errors are approximately \( \pm 1000 \) K in temperature and \( \pm 0.3 \) dex in \( \log g \), and the results are listed in Table 1.

We estimated the projected rotational velocity \( V \sin i \) by comparing the observed He i and He ii line widths with those found by convolving the model profiles from Lanz & Hubeny (2003) with instrumental and rotational broadening functions. The instrumental broadening of our spectra was represented by a Gaussian with FWHM = 52 km s\(^{-1}\). The rotational broadening function (Gray 1992) was computed assuming a spherical star with a linear limb-darkening coefficient determined from Wade & Rucinski (1985). A grid of theoretical profiles was computed for a range of \( V \sin i \), and a \( \chi^2 \) minimization was used to determine the best fit to the data. The results and errors are listed in Table 1. Note that this approach does not account for variations in the intensity profile across the disk, nor the nonspherical shape and gravity darkening found in rapid rotators (Townsend et al. 2004). These shortcomings are most important for the two rapid rotators in our sample, BD +60 513 and BD +60 594, and we have probably underestimated the true rotational broadening in these cases.

5. RADII FROM FITS OF THE SPECTRAL ENERGY DISTRIBUTIONS

All of our targets are considered to be members of the Cas OB6 association based on their position in the sky and placement in the Hertzsprung-Russell diagram (HRD; Voroshilov et al. 1985; Garmany & Stencel 1992; Walborn 2002). The proper-motion study of IC 1805 made by Vasilevsksis et al. (1965) confirms membership for five targets but curiously indicates a somewhat discrepant proper motion for HD 15558. We assume that all the targets share a common distance, so their observed fluxes are directly related to their relative stellar radii. Since we have determined temperature and multiplicity for each target, we have all parameters needed to estimate their radii as a function of distance by comparing their observed and model flux distributions, and we can then check the distance by comparing the derived and model radii (as a function of temperature). The angular diameter of the limb-darkened disk \( \theta_{\text{LD}} \) (in units of radians) is found by the inverse square law:

\[
\frac{f_i(\text{observed})}{F_i(\text{emitted})} 10^{-0.4i} = \left(\frac{R_i}{d}\right)^2 = \frac{1}{4} \theta_{\text{LD}}^2,
\]

where the ratio of the observed and emitted fluxes (reduced by the effects of interstellar extinction \( A_i \)) depends on the square of the ratio of stellar radius \( R_i \) to distance \( d \). All our targets have a large reddening, so we need to fit the relation above over the entire observed spectrum in order to account reliably for the wavelength-dependent shape of the extinction curve. We adopted the extinction curve law from Fitzpatrick (1999) that is a function of the reddening \( E(B-V) \) and the ratio of total to selective extinction \( R = A_V/E(B-V) \). We then fitted the observed flux distribution of each target to determine \( E(B-V) \), \( R \), and angular diameter \( \theta_{\text{LD}} \).

We compiled data on the observed fluxes from the UV to the near-IR. There are low-dispersion International Ultraviolet Explorer (IUE) spectra available (SWP and LWR cameras) for all of the targets except DN Cas, HD 17520, HD 237019, and BD +60 586, and we formed binned versions of the UV spectral fluxes from these. We found Johnson \( UBV \) photometry for all the targets (plus Cousins \( RI \) in a few cases), and these magnitudes were transformed to fluxes using the calibration of Colina et al. (1996). We also included Strömgren photometry for 7 of the 13 targets using the calibration of Gray (1998). All of the stars are included in the Two Micron All Sky Survey (2MASS) All-Sky Catalog of Point Sources (Cutri et al. 2003), and we converted the \( JHK \) magnitudes to fluxes using the calibration of Cohen et al. (2003).

The model fluxes were selected from the grid of Lanz & Hubeny (2003) based on the values of \( T_{\text{eff}} \) and \( \log g \) in Table 1. These model fluxes are based on plane-parallel, non-LTE, line-blanketed atmospheres that do not include the effects of stellar wind mass loss. However, the inclusion of winds generally alters...
of the spectral energy distributions only at wavelengths much larger and much smaller than those considered here (Martins et al. 2005), so wind effects are insignificant in this context. We co-added model fluxes using the observed flux ratios for the spectroscopic binaries and using the observed magnitude difference for close visual binaries (details are given in next section for individual stars). Note that in the two single-lined binary cases, fit was based on the optical fluxes alone in this case.

The stellar radius is found by

\[ R_* = \frac{(10.43 R_\odot) \theta_{LD} d}{\mu} \]                      

We compared the observed and model flux distributions and made a grid search fit to find the values of \( E(B - V) \) and \( R \) that produced the lowest \( \chi^2 \) residuals (by weighting the observations approximately equally between the UV, optical, and near-IR bands). The normalization term in the final fit yields the angular diameter. Our results are gathered in Table 3, which lists the star name, home cluster, angular diameter (in microarcseconds), derived radius (see below), \( E(B - V) \), and \( R \). The accompanying error estimates for \( E(B - V) \) and \( R \) are derived from the point in grid search where \( \chi^2 \) attains a value higher by 1 than the normalized minimum value, and the errors in diameter are based on the scatter in the flux distribution fit and the errors propagated from uncertainty in \( T_{eff} \), \( E(B - V) \), and \( R \). The model fits were generally excellent, but they were less satisfactory for the spectra of the three hottest stars, where it was difficult to find a consistent match of the far- and near-UV fluxes. We also found that the \( JHK \) fluxes were too large relative to the optical fluxes for HD 17520, which we suspect is due to an IR excess associated with the Be disk emission of the companion (Porter & Rivinius 2003), so the fit was based on the optical fluxes alone in this case.

**Table 3**

| Object Name     | Cluster Membership | \( \theta_{LD} \) (arcsec) | \( R_\odot^* \) (\( R_\odot \)) | \( E(B - V) \) (mag) | \( R \) [\( A_V/E(B - V) \)] |
|-----------------|--------------------|-----------------------------|---------------------------------|---------------------|-----------------------------|
| DN Cas A        | IC 1805            | 36 (1)                      | 7.4 (7)                         | 0.97 (1)            | 3.13 (1)                    |
| DN Cas B        | IC 1805            | 26 (3)                      | 5.3 (8)                         | 0.97 (1)            | 3.13 (1)                    |
| BD +60 497A     | IC 1805            | 49 (4)                      | 10.0 (13)                       | 0.89 (2)            | 2.94 (6)                    |
| BD +60 497B     | IC 1805            | 32 (5)                      | 6.5 (11)                        | 0.89 (2)            | 2.94 (6)                    |
| BD +60 501      | IC 1805            | 36 (3)                      | 7.2 (9)                         | 0.76 (2)            | 3.12 (5)                    |
| HD 15558A       | IC 1805            | 81 (10)                     | 16.4 (26)                       | 0.84 (3)            | 3.07 (9)                    |
| HD 15558B       | IC 1805            | 28 (4)                      | 5.6 (9)                         | 0.84 (3)            | 3.07 (9)                    |
| HD 15570        | IC 1805            | 95 (16)                     | 19.3 (37)                       | 1.01 (2)            | 3.13 (7)                    |
| HD 15629        | IC 1805            | 57 (7)                      | 11.6 (19)                       | 0.76 (3)            | 3.09 (12)                   |
| BD +60 513      | IC 1805            | 43 (4)                      | 8.7 (12)                        | 0.83 (3)            | 3.00 (5)                    |
| BD +62 424      | IC 1805            | 49 (3)                      | 9.9 (11)                        | 0.75 (2)            | 3.02 (4)                    |
| HD 17505Ab      | IC 1848            | 56 (6)                      | 11.4 (16)                       | 0.72 (3)            | 2.83 (9)                    |
| HD 17505Aa1     | IC 1848            | 49 (8)                      | 10.0 (18)                       | 0.72 (3)            | 2.83 (9)                    |
| HD 17505Aa2     | IC 1848            | 49 (3)                      | 8.9 (15)                        | 0.72 (3)            | 2.83 (9)                    |
| HD 17505B       | IC 1848            | 38 (11)                     | 7.7 (24)                        | 0.57 (5)            | 3.10 (120)                  |
| HD 17520Ab      | IC 1848            | 14 (4)                      | 2.8 (9)                         | 0.57 (5)            | 3.10 (120)                  |
| HD 17520B       | IC 1848            | 34 (10)                     | 7.0 (22)                        | 0.57 (5)            | 3.10 (120)                  |
| HD 17520C       | IC 1848            | 14 (4)                      | 2.8 (9)                         | 0.57 (5)            | 3.10 (120)                  |
| HD 237019       | IC 1848            | 37 (1)                      | 7.4 (17)                        | 0.77 (1)            | 3.20 (6)                    |
| BD +60 586A     | IC 1848            | 45 (3)                      | 9.1 (11)                        | 0.58 (3)            | 3.27 (32)                   |
| BD +60 586B     | IC 1848            | 25 (2)                      | 5.1 (7)                         | 0.58 (3)            | 3.27 (32)                   |
| BD +60 594A     | IC 1848            | 35 (4)                      | 7.2 (11)                        | 0.68 (2)            | 2.88 (10)                   |
| BD +60 594B     | IC 1848            | 10 (2)                      | 2.0 (5)                         | 0.68 (2)            | 2.88 (10)                   |

* Radius derived for a distance of 1.886 kpc.
stellar radii are directly proportional to the assumed distance, and, for example, if we adopted the larger distance $d = 2.34$ kpc advocated by Massey et al. (1995), then the stars would be approximately 24% larger, a size that would be difficult to reconcile with the evolutionary models of Lejeune & Schaerer (2001). We show below that our derived radii for DN Cas agree well with those determined independently from an analysis of the eclipsing light curve.

6. RESULTS FOR INDIVIDUAL OBJECTS

6.1. Stars with Constant Radial Velocity

**BD +60 501.**—The two Balmer lines H$\gamma$ and H$\beta$ have the lowest measured velocities in the template spectrum, so they were omitted from the average (i.e., based only on the He i and He ii lines). Our observed mean velocity of $-57.9 \pm 1.1$ km s$^{-1}$ is reasonably consistent with the mean of $-49.9 \pm 2.5$ km s$^{-1}$ from the work of Rauw & De Becker (2004), who also find the star to be a constant-velocity object.

**HD 15558.**—There is evidence here for a velocity progression (lower velocity among low-excitation lines), so we used only the lines He ii $\lambda\lambda 4199, 4541$ to set the radial velocity of the template spectrum. There are no obvious variations over the five-night run, but the star may be a long-period, single-lined binary (Garmany & Massey 1981). According to the recent orbital elements from Stickland & Lloyd (2001), the orbital period is 439.64 $\pm$ 0.26 days and the predicted velocity at the mean time of our observations is $-21.8$ km s$^{-1}$ (near the peak of the radial velocity curve), which is $27.7$ km s$^{-1}$ higher than our observed mean of $-49.5 \pm 2.9$ km s$^{-1}$. This discrepancy could be caused by differences in measurement techniques or errors in the preliminary orbital solution of Stickland & Lloyd (2001). A long-term study by M. De Becker & G. Rauw (2005, private communication) has recently confirmed the binary status of this star and will provide an improved orbital solution. Therefore, we list this star as “constant (SB1)” in Table 1, denoting both a lack of short-term line variability in our observations and the recent confirmation of binarity. There is a visual companion at a separation of 9$^\prime\prime$ that is 2.6 mag fainter (Perryman 1997).

**HD 15570.**—There is a clear velocity progression between spectral lines with the low-excitation lines (H, He i) appearing at a much lower velocity than the high-excitation lines (He ii, N v). There is also a correlation with line width in the sense that stronger, broader lines have lower velocity. We adopted a radial velocity for the template spectrum from the lines He ii $\lambda\lambda 4199$ and N v $\lambda\lambda 4604, 4620$. The night-to-night scatter between velocities is small except on the last night when the measured velocity is $\approx 10$ km s$^{-1}$ more positive than the mean from the other nights. Curiously, the O emission lines do not show any shift in velocity on the final night. We suspect that the change was related to an atmospheric fluctuation rather than orbital motion. Underhill & Gilroy (1990) also consider this star to be a velocity variable. De Becker et al. (2005) find no evidence of orbital motion and discuss rotational modulation as a possible source of line variability.

**HD 15629.**—The low-excitation Balmer and He i lines have systematically low velocities, so the template mean was adopted from the lines He ii $\lambda\lambda 4199, 4541, 4686$. We find no evidence of velocity variability. Underhill & Gilroy (1990), on the other hand, flag this star as possibly variable, but given the presence of line-to-line velocity differences, some of the velocity differences found by various authors are probably due to the line sample adopted (and not to orbital motion). De Becker et al. (2005) also find no evidence of orbital motion for this star.

**BD +60 513.**—This star has a broad-lined spectrum, and we selected only the strongest lines to measure the template mean velocity (H$\gamma$, H$\beta$, He i $\lambda\lambda 4471, 4921$, He ii $\lambda\lambda 4541, 4686$). We confirm the lack of short-term velocity variability reported by Rauw & De Becker (2004), and our mean velocity of $-58.6 \pm 2.8$ km s$^{-1}$ is probably consistent with their mean value, $-44.3 \pm 10.5$ km s$^{-1}$.

**BD +62 424.**—The H$\beta$ line had the lowest velocity measured, so we omitted the H Balmer lines and weak lines in forming the template mean velocity. The star appears to be radial velocity constant over the run, and our mean velocity, $-34.3 \pm 1.7$ km s$^{-1}$, is probably consistent with that from Abt et al. (1972) of $-43.5 \pm 6.3$ km s$^{-1}$.

**HD 237019.**—All the strong lines were used to form the template mean velocity. The star has a constant radial velocity over the run. Our mean velocity of $-43.4 \pm 1.1$ km s$^{-1}$ is higher than the only other value in the literature of $-68$ km s$^{-1}$ from Seyfert & Popper (1941).

**BD +60 586.**—We omitted H$\beta$ and He i $\lambda 4921$ from the template mean calculation, since both lines had significantly lower velocities. The star was radial velocity constant during the run. Our mean velocity, $-48.1 \pm 1.0$ km s$^{-1}$, agrees with that found by Abt et al. (1972), $-42.3 \pm 7.1$ km s$^{-1}$, but is lower than the value reported by Conti et al. (1977), $-27.9 \pm 1.2$ km s$^{-1}$. The star is a member of a visual multiple system (ADS 2194) with companions at separations of 7$''$ (2$\Delta V = 1.7$) and 13$''$ (2$\Delta V = 4.7$).

6.2. Single-Lined Spectroscopic Binaries

**HD 17520.**—The H$\beta$ line has a blueshifted emission component that also appears weakly in H$\gamma$. The Balmer lines and weak lines were omitted in the calculation of the template mean velocity. The velocities show a progressive increase over the duration of the run, reaching values similar to those observed by Conti et al. (1977) and Liu et al. (1989). Walter (1992) showed how H$\alpha$ emission developed between 1985 and 1991 leading to a spectral classification as a Be-type star (probably its status in 2003 as well since the emission is so apparent in H$\beta$). This star is also the central object in a visual multiple system (ADS 2165) with a close companion at a separation of 0$''$.335 and with a magnitude difference of 0.52 mag (Perryman 1997; Walborn 2002).
spectroscopic observations of DN Cas have appeared in the literature, although Davidson (1980) claimed that such observations were underway. Additionally, the only photometry that has been published, other than in phase-folded figures, is the compromised data of Frazier & Hall (1974). This overlooked opportunity to obtain masses for this massive binary is rectified here, where we present the first full orbital solution based on our spectroscopy and photometry from the literature.

We determined velocities for both components using non-linear least-squares fits of two templates with temperatures and rotational velocities given in Table 1 and with a monochromatic flux ratio in the blue of $F_2/F_1 = 0.41 \pm 0.10$. Our final radial velocities are given in Table 4, and the spectroscopic orbital elements are given in Table 5. The period was adopted from the

![Graph](image_url)
Here we adopt the time of minimum light (primary superior conjunction) to set the zero point of orbital phase and epoch. Eclipses were observed monochromatic photometry data folded on the given ephemeris and plotted with the Wilson-Devinney code light curve.

eclipse timing observations collected by Kreiner et al. (2001). Here we adopt the time of minimum light (primary superior conjunction) to set the zero point of orbital phase and epoch \( T_{\text{min}} \), and our derived epoch agrees with the current photometric epoch of minimum light from Kreiner et al. (2001) of HJD 2,452,501.928 \( \pm 0.007 \). Eccentric fits offered no significant improvement over circular fits (Lucy & Sweeney 1971), so circular elements are given in Table 5. The radial velocities and orbital solutions are shown in Figure 4.

The only published photometric data of DN Cas are those from Frazier & Hall (1974), who made \( UBV \) photoelectric observations between 1971 October and 1972 December. Unfortunately, their comparison star, HD 14817, is itself an eclipsing binary as mentioned above. However, the light curve of HD 14817 from \( \text{Hipparcos} \) (Perryman 1997) is flat outside of eclipse, so we can retain the results from Frazier & Hall (1974) that are based on outside-of-eclipse times for HD 14817. These measurements are plotted against orbital phase in Figure 5, and because there remain significant phase gaps in this restricted set, our light-curve fitting results should be considered preliminary.

We calculated a model light curve using the Wilson-Devinney code (Wilson & Devinney 1971; Wilson 1990). Most of the model parameters are set by the spectroscopic orbital solution (Table 5), the adopted temperatures and gravities from the spectra (Table 1), the associated limb-darkening parameters (Van Hamme 1993), and the monochromatic flux ratio (which sets the ratio of the mean stellar radii). In addition, we set the bolometric albedo and gravity darkening exponents to 1.0, which are valid approximations for stars with fully radiative envelopes. We then fitted the observed \( V \)-band observations by varying the inclination \( i \) and the stellar radii. The ratio of stellar radii was set according to the temperatures and gravities in Table 1,

\[
\frac{R_2}{R_1} = \left\{ \frac{f_i(\text{observed})_2}{f_i(\text{observed})_1} \right\} \frac{\left[ F_v(\text{emitted})_2 \right]}{\left[ F_v(\text{emitted})_1 \right]}^{0.5}
\]

\[
= 0.72 \pm 0.09.
\]

If the stars are rotating synchronously, then the ratio of projected rotational velocities also yields the ratio of radii as

\[
\frac{R_2}{R_1} = \frac{(V \sin i)_2}{(V \sin i)_1} = 0.72 \pm 0.15,
\]

and the good agreement between these two estimates suggests that synchronous rotation has been attained in this close binary. We found that the radii obtained from a fit of the spectral energy distribution and distance of 1.9 kpc (Table 3) led to an excellent fit of the light curve for an orbital inclination of \( i = 79^\circ 2 \pm 1^\circ 1 \) (Fig. 5), and the derived stellar masses are \( M_1 = 19.2 \pm 0.8 M_\odot \) for the O8 V primary and \( M_2 = 13.9 \pm 0.5 M_\odot \) for the B0.2 V secondary (Table 5). The mass of the primary is consistent with its spectral type for the calibration given by Martins et al. (2005), and the mass of the secondary agrees well with the mass for its spectral type according to the relations given by Hanson et al. (1997). Both stars are well within their respective Roche lobes.

There is a companion to DN Cas in the 2MASS catalog at a separation of 5.7\( \arcsec \) and at a position angle of 89\( ^\circ \). However, inspection of the 2MASS images at this position suggests that this object is probably not a real star since its point-spread function looks very sharp and unlike those of nearby stars. Furthermore, we see no evidence of such a tertiary in our observations. Since the near east-west alignment of the putative object and DN Cas is the same as the slit orientation for our KPNO 4 m spectra, the tertiary should appear in our spectra with a separation from DN Cas of 7.7 pixels in the spatial dimension. The average 2MASS magnitude difference is 2.13 mag (for \( H, K \)), so that the flux ratio of tertiary to binary is 0.14 in the near-IR. However, our blue spectra show no sign of a companion brighter than \( \Delta m = 5 \), so if this companion is real, it is probably a very red foreground star that can be ignored in our analysis. We found that the introduction of a third light contribution of 14% in the \( V \)-band light curve results in an increase in the inclination by 3\( ^\circ \), but since the actual flux contribution of a tertiary is probably much smaller than this, any changes in our results due to the presence of a tertiary are probably comparable to or less than the errors quoted in Table 5.

6.4. The Double-Lined System BD +60 497

The spectra of BD +60 497 show clear evidence that this star is a double-lined spectroscopic binary. In their recent paper, Rauw & De Becker (2004) give the first orbital solution for BD +60 497, and they find an orbital period of 3.96 days. We measured velocities using the scheme outlined in §3 by matching template synthetic profiles for models with temperatures and gravities given in Table 1, together with an estimated flux ratio of \( F_2/F_1 = 0.35 \pm 0.08 \). We note, however, that in some of the lines (for example, He I \( \lambda 4387, 4471 \)) we found that the approaching, blueshifted component (primary and secondary) often appeared to be weaker than expected based on this fixed flux ratio (Rauw & De Becker [2004] also observed a similar variation). We doubt that this variation has a significant effect on our measurements, which appear in Table 6.

We combined our velocities with those of Rauw & De Becker (2004) in making the orbital solution. A number of alias periods remain viable, but we obtained the best fit with a period close to that obtained by Rauw & De Becker (2004) (see Table 7). However, unlike Rauw & De Becker (2004), we found that a significant eccentricity is present, and we made full elliptical fits for both components. Otherwise, our results are in reasonable agreement with their solution (compared in Table 7). The velocities and orbital fit are shown folded on the derived ephemeris in

[Fig. 5.—DN Cas \( V \)-band photometry data folded on the given ephemeris and plotted with the Wilson-Devinney code light curve.]

\[ R_2/R_1 = \frac{(V \sin i)_2}{(V \sin i)_1} = 0.72 \pm 0.15, \]

and the good agreement between these two estimates suggests that synchronous rotation has been attained in this close binary. We found that the radii obtained from a fit of the spectral energy distribution and distance of 1.9 kpc (Table 3) led to an excellent fit of the light curve for an orbital inclination of \( i = 79^\circ 2 \pm 1^\circ 1 \) (Fig. 5), and the derived stellar masses are \( M_1 = 19.2 \pm 0.8 M_\odot \) for the O8 V primary and \( M_2 = 13.9 \pm 0.5 M_\odot \) for the B0.2 V secondary (Table 5). The mass of the primary is consistent with its spectral type for the calibration given by Martins et al. (2005), and the mass of the secondary agrees well with the mass for its spectral type according to the relations given by Hanson et al. (1997). Both stars are well within their respective Roche lobes.

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Figure 6. Assuming masses typical of their spectral classes (Martins et al. 2005), the system probably has an inclination of $i \approx 47^\circ$, so we do not expect eclipses to occur.

6.5. The Multiple-Star System HD 17505

The star system HD 17505 dominates the center of the cluster IC 1848, and there are seven visual companions noted by Mason et al. (1998) at distances ranging from 2.71 to 12.40. Multiplicity in the central object HD 17505A has been discussed by Conti & Aachslcher (1971), Walborn (1973), Howarth et al. (1997), and Walborn (2002), and it is known to contain at least three stars since the single IUE spectrum shows evidence of three Doppler-shifted components (Stickland & Lloyd 2001). Walborn (2002) pointed out that the multiplicity of the star is the reason it appears so bright in the HRD. We obtained 10 spectra of HD 17505A over the five-night duration of the KPNO 4 m run, and we found that the lines were strongly blended in all but a few spectra corresponding to the maximum Doppler separation of the binary motion. However, we also obtained an additional 32 red spectra of the target in 2000 with the KPNO 0.9 m coude’ feed telescope (§3), and these turned out to be particularly helpful in finding the orbital elements.

Initial inspection of these higher dispersion red spectra showed that all three components were cleanly resolved in the He i profiles observed near orbital quadrature phases and indicated an orbital period of approximately 8 days. We measured radial velocities by first making unconstrained three-Gaussian fits of the least blended feature, He i 7065, in the quadrature phase spectra. We obtained a preliminary orbital solution from this subset of spectra for the two moving components, and then we used this solution to estimate the Doppler shifts for all the coude’ feed observations. Then we made three-component Gaussian fits of He i 7065 in all the red spectra using these starting estimates for the Doppler shifts and constraining the widths to be the mean values found from the fits of the quadrature phase spectra. A new orbital solution was obtained, and the solution was used to make a three-component reconstruction of the individual component spectra using the Doppler tomography algorithm (Bagnuolo et al. 1994). We subtracted the reconstructed spectrum of the stationary component from each spectrum and renormalized them to end up with a set corresponding to the spectra of the double-lined binary alone. We then measured the radial velocities using the nonlinear, least-squares method by matching template spectra to the profiles of H$\alpha$, He i 6678/He ii 6683, and He i 7065. The template spectra for the two components were taken from their respective tomographic reconstructions ($F_2/F_1 = 1.0 \pm 0.1$) rather than using the model profiles from Lanz & Hubeny (2003) because we were concerned that some of these red features are not well represented.
was the lines were too badly blended for reliable measurement. Our conjunction phase observation (from HJD 2,451,892.743) where the synthetic models. This approach worked well except in one

\begin{tabular}{lcccc}
HJD & Orbital Phase & $V_1$ & $(O-C)_1$ & $V_2$ & $(O-C)_2$ \\
(-2,450,000) & & (km s\(^{-1}\)) & (km s\(^{-1}\)) & (km s\(^{-1}\)) & (km s\(^{-1}\)) \\
1817.887 & 0.772 & -181.1 & 5.3 & 143.7 & 5.2 \\
1818.890 & 0.889 & -184.6 & -2.8 & 136.3 & 2.6 \\
1819.873 & 0.004 & -77.9 & -1.3 & 18.9 & -6.9 \\
1820.887 & 0.122 & 63.4 & -0.2 & -124.7 & -6.7 \\
1821.864 & 0.236 & 148.2 & 15.8 & -190.4 & -1.9 \\
1822.865 & 0.353 & 120.9 & 3.7 & -171.9 & 1.1 \\
1823.814 & 0.464 & 49.2 & 0.0 & -105.9 & -2.9 \\
1823.924 & 0.476 & 35.7 & -3.4 & -95.7 & -2.8 \\
1824.825 & 0.582 & -45.5 & 4.8 & 1.0 & 2.0 \\
1824.950 & 0.596 & -54.8 & 8.4 & 16.6 & 4.5 \\
1830.842 & 0.284 & 142.9 & 7.5 & -187.7 & 3.8 \\
1830.946 & 0.296 & 136.5 & 2.5 & -180.0 & 10.1 \\
1888.797 & 0.045 & -29.1 & -4.5 & -34.0 & -6.5 \\
1889.746 & 0.156 & 91.2 & -2.1 & -148.5 & -0.0 \\
1892.743 & 0.506 & ... & ... & ... & ... \\
1893.774 & 0.626 & -77.2 & 11.8 & 25.5 & -13.0 \\
1894.771 & 0.742 & -168.5 & 4.5 & 131.3 & 6.6 \\
1894.845 & 0.751 & -183.3 & -6.0 & 136.3 & 7.2 \\
1895.721 & 0.853 & -196.3 & -1.8 & 144.7 & -2.1 \\
1895.811 & 0.864 & -184.3 & 7.5 & 144.7 & 0.6 \\
1896.651 & 0.962 & -132.7 & 8.0 & 76.7 & 1.4 \\
1896.764 & 0.977 & -103.8 & 4.1 & 45.3 & -12.6 \\
1897.652 & 0.078 & 6.0 & -10.1 & -71.1 & -1.8 \\
1897.784 & 0.094 & 21.0 & -12.9 & -87.4 & 0.1 \\
1898.656 & 0.196 & 114.5 & -3.8 & -177.2 & -3.2 \\
1898.793 & 0.212 & 126.0 & 0.8 & -173.9 & 7.2 \\
1899.658 & 0.312 & 133.2 & 2.5 & -182.4 & 4.3 \\
1899.791 & 0.328 & 124.9 & -1.5 & -182.2 & 0.2 \\
1900.653 & 0.429 & 69.7 & -4.9 & -136.9 & -7.7 \\
1900.784 & 0.444 & 60.6 & -3.2 & -124.9 & -6.7 \\
1901.637 & 0.543 & -32.7 & -15.8 & -141.4 & 21.2 \\
1901.770 & 0.559 & -27.6 & 3.0 & -24.4 & -3.0 \\
\end{tabular}

by the synthetic models. This approach worked well except in one conjunction phase observation (from HJD 2,451,892.743) where the lines were too badly blended for reliable measurement. Our final radial velocities from the red spectra for the close binary are collected in Table 8.

The estimate for the orbital period from the red spectra alone was $P = 8.572 \pm 0.003$ days, and the error was sufficiently small that we could determine the cycle number and approximate orbital phase corresponding to the times of the KPNO 4 m observations in 2003 October. We measured the line width of the He i 4471 feature in these blue spectra to find that maximum velocity separation occurred at HJD 2,452,915.469 $\pm$ 0.12, and we compared this to the time of the same orbital phase for the 2000 coude' observations to arrive at a revised period estimate of $P = 8.5710 \pm 0.0008$ days. The final orbital solution is given in Table 9, and the radial velocity curve appears in Figure 7. The predictions from this fit for the time of the IUE observation are +128 and $-184$ km s\(^{-1}\), in reasonable agreement with the observed velocities of +142 and $-165$ km s\(^{-1}\) (Stickland & Lloyd 2001), but we did not include the IUE measurements in the final solutions because of possible systematic differences between velocities measured in the UV and optical.

We made tomographic reconstructions of both the blue and red spectra using the orbital solution in Table 9 and assuming flux contributions of 30%, 30%, and 40% for Aa1, Aa2, and Ab,

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respectively (based on the relative equivalent widths of the \( \text{He} \, \lambda 6678 + \text{He} \, \lambda 6683 \) feature). Both the blue and red spectral reconstructions are plotted in Figure 8. We determined radial velocities of the stationary component Ab of \(-30 \pm 3\) and \(-38 \pm 5\) from the blue and red reconstructed spectra, respectively. Once again, the optical velocities are systematically more negative than the UV-determined value for Ab of \(-13.5 \, \text{km s}^{-1}\) (Stickland & Lloyd 2001). Our results are summarized in a schematic representation of the system hierarchy in Figure 9. Note that the coude' feed spectra were made with a \( 2.5 \) wide slit in an approximately north-south orientation, so the B component (at a separation of \( 2.15\) and a position angle of \( 93^\circ \) with a magnitude difference of \( \Delta H_p = 1.66 \) mag; Perryman 1997) may be slightly mixed in with the light from the Ab component for the red reconstructed spectrum. However, the 4 m spectra were made with an east-west slit orientation, so it was possible to separate the A and B components spatially along the slit. A single extracted spectrum of the B component is shown in Figure 8 along with the tomographic reconstructions of the Aa1, Aa2, and Ab components. The spectrum of the B component is clearly different and indicates a cooler effective temperature than the stars of component A; thus, the spectroscopic triple we detect is clearly related to A alone. The hottest and brightest of all the components is Ab, and its spectral type agrees with that obtained by Walborn (1972) for the entire system.

The stars of the Aa1-Aa2 central binary are far enough apart that no eclipses should be seen, and indeed none are observed (Perryman 1997). For example, if we adopt the masses and radii associated with O-type main-sequence stars from Martins et al. (2005), then the expected system inclination should be \( i = 65^\circ \) (based on the \( M \sin i^2 \) values in Table 9), while eclipses would be observed only if \( i > 73^\circ \) (based on the semimajor axis in Table 9). The orbital period of B around the triple A can be estimated as \( \approx 27,000 \) yr if we take the projected angular separation as the semimajor axis and the distance is \( 1.9 \) kpc. The orbital period of the Aa-Ab central triple is probably much less. The system was unresolved in the speckle survey of Mason et al. (1998), and if we take their resolution limit of \( 35 \) mas as the projected semimajor axis, then the orbital period should be \(< 61 \) yr. The target is clearly worthy of continued spectroscopic and interferometric observation to search for evidence of such orbital motion.

### 7. SUMMARY

We have monitored 13 O stars in the Cas OB6 association in search of multiple-star systems. Five of the 13 were found to have radial velocity variations related to orbital motion. Of these five, we have determined orbital solutions for three double-lined systems: the eclipsing binary DN Cas, BD +60 497, and the spectroscopic triple HD 17505A. The remaining two binaries, BD +60 594 and HD 17520, are single-lined binaries with periods longer than 5 days. In addition to these five, one of our targets, HD 15558, is probably a long-period binary (Garmany & Massey 1981; Stickland & Lloyd 2001). One other bright O star in Cas OB6 is HD 16429, which was found by McSwain (2003) to be a spectroscopic triple. If we include all of these objects in the known sample of observed O stars in the Cas OB6 association, then we find that 7 of the 14, or 50%, are spectroscopic binaries. This value is well within the range found by Mason et al. (1998) for spectroscopic binaries among stars in clusters and associations.

We note that two systems, HD 17505 and DN Cas, are of special interest and would benefit from follow-up observations. We are planning to obtain new three-color photometry of the eclipsing binary DN Cas in order to make a better fit of the light curve and to determine an accurate mass of the O8 V primary star. The multiple system HD 17505 is a fascinating system with at least four O stars (with a total mass close to \( 100 M_\odot \)) that are apparently gravitationally bound. High-resolution studies via interferometry to separate the Aa and Ab components would produce a more accurate kinematic model for the system and would help us better understand stellar interactions in these systems and their role in the formation and ejection of massive stars.

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