Yang-Yang thermometry and momentum distribution of a trapped one-dimensional Bose gas

M. J. Davis,1 P. B. Blakie,2 A. H. van Amerongen,3 N. J. van Druten,3 and K. V. Kheruntsyan†

1The University of Queensland, School of Mathematics and Physics, Qld 4072, Australia
2Jack Dodd Centre for Quantum Technology, Department of Physics, University of Otago, Dunedin, New Zealand
3Van der Waals-Zeeman Institute, University of Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands

(Dated: November 8, 2018)

We describe the use of the exact Yang-Yang solutions for the one-dimensional Bose gas to enable accurate kinetic-energy thermometry based on the root-mean-square width of an experimentally measured momentum distribution. Furthermore, we use the stochastic projected Gross-Pitaevskii theory to provide the first quantitative description of the full momentum distribution measurements of Van Amerongen et al., Phys. Rev. Lett. 100, 090402 (2008). We find the fitted temperatures from the stochastic projected Gross-Pitaevskii approach are in excellent agreement with those determined by Yang-Yang kinetic-energy thermometry.

PACS numbers: 05.30.Jp, 03.75.Hh, 05.70.Ce

Ultracold gases offer a unique opportunity to study fundamental problems in quantum many-body physics, allowing experimental observations to be compared directly with microscopic theories. An area of significant recent interest has been the measurement of thermodynamic relations [1–4]. The one-dimensional (1D) Bose gas with repulsive interactions has emerged as a paradigm system because exact solutions are available for both eigenstates [4] and thermodynamic quantities [5] (see Ref. [6] and references therein). Furthermore, this system exhibits a surprisingly rich variety of regimes [7, 8] connected by broad crossovers. Most studies of the 1D Bose gas have focused on the position-space distributions [3, 8–10] and local correlations [7, 8, 11–14], which can be directly obtained from the exact theories.

A recent experiment by Van Amerongen et al. [10] measured the position and momentum distributions of a trapped 1D Bose gas throughout the crossover from an ideal gas to the quasicondensate regime. The position-space measurements were compared with the Yang-Yang (YY) thermodynamic solutions [5] within the local density approximation (LDA), and showed smooth behavior throughout the crossover. In contrast, the momentum distributions showed a pronounced temperature dependence, and have been unexplained by theory to date. Previous work on the momentum properties of the 1D Bose gas has focused on limiting cases [12, 13, 15].

Here we investigate the momentum properties of the 1D Bose gas and their application to thermometry through measurements of the system kinetic energy. Our methods provide a reliable foundation for accurate thermometry in all regimes of a 1D Bose gas with repulsive interactions, including the strongly-correlated regime. This approach is reminiscent of molecular dynamics calculations, where the average kinetic energy per particle is a direct measure of the temperature [16].

First, we use the exact YY thermodynamic formalism [5] to calculate the root-mean-square (rms) width of the momentum distribution, which is equivalent to determining the average kinetic energy per particle. In combination with the LDA for trapped (nonuniform) quasi-1D systems, we show how the YY kinetic energy results can be applied to accurate thermometry for a broad range of conditions that are relevant to ongoing experimental and theoretical efforts. We refer to this approach as YY thermometry. Second, we present the first quantitative calculation of the full momentum distribution for a trapped quasi-1D Bose gas using the stochastic projected Gross-Pitaevskii technique [17]. (Calculations of the position-space distribution based on a related formalism [18] were recently described in Ref. [19].) We find excellent agreement between the SPGPE results and the momentum-space measurements reported in Ref. [10]. Finally, we compare the kinetic energy and temperature predictions of the YY thermometry with those of the SPGPE and also find excellent agreement between the two approaches. While the SPGPE technique is limited to the degenerate yet high-temperature weakly interacting regime, the YY thermodynamic formalism applies to all repulsive interaction strengths and temperatures and therefore its usefulness extends beyond the regimes studied here.

A uniform 1D Bose gas in the thermodynamic limit is completely characterized by two parameters [4–9]: the dimensionless interaction strength $\gamma = mg^2/\hbar^2 \rho$ and the reduced temperature $t \equiv 2k_B T \hbar^2 / m g^2$, where $\rho$ is the linear (1D) density, $g \simeq 2\hbar \omega_\perp a$ is the effective 1D coupling strength [20], $a$ is the 3D $s$-wave scattering length, and $\omega_\perp$ the transverse radial harmonic trapping frequency [21]. The 1D regime is realized when the transverse excitation energy $\hbar \omega_\perp$ is much larger than the thermal energy $k_B T$ and chemical potential $\mu$. While the YY thermodynamic equations do not directly yield the 1D momentum distribution, $n(k_z)$, here we show how they can be used to obtain the average kinetic energy per particle

$$E_{\text{kin}}/N = \hbar^2 \langle k_z^2 \rangle/2m,$$

where $\langle k_z^2 \rangle^{1/2} = \left[ \int \! dk_z \, k_z^2 \, n(k_z) / N \right]^{1/2}$ is the rms width of the distribution $n(k_z)$, and $N = \int \! dk_z \, n(k_z)$ is the total atom number.

Solutions to the YY thermodynamic equations yield a unique value for the total energy per particle, $E/N$ [5], for a given temperature and interaction strength. Using the Helmann-Feynman theorem, the YY solutions can also be used to determine the local pair correlation function
parameter $\mu/k$ of the ground transverse mode is $\langle \rho \rangle_{\text{YY}}$, the system becomes fermionized, leading to an increase in $E_{\text{kin}}/N$ for both negative and positive $\mu$. For $t \ll 1$, and $1 \ll \gamma \ll t^{-1/2}$ in the Tonks-Girardeau regime, $E_{\text{kin}}/Nk_B T$ becomes larger than the equipartition value of $1/2$.

In order to apply the YY results to trapped (nonuniform) quasi-1D systems as realized in experiment [10], we utilise the LDA. The trapping potential is $U(r) = V(z) + \frac{1}{2} m \omega_z^2 z^2$, where $V(z) \geq \frac{1}{2} m \omega_1^2 z^2$ is the potential in the weakly confined longitudinal direction with $\omega_z \ll \omega_1$ [23]. We calculate the particle number density $\rho(z)$ by treating the trapped system as a collection of sufficiently small uniform systems of length $\Delta z$, with the local chemical potential $\mu(z) = \mu - V(z)$ [8,9] where $\mu$ is the global chemical potential. The density of the trapped system in the ground transverse mode, $\rho_0(z)$, is evaluated as $\rho_0(z) = \rho_{\text{YY}}[\mu(z), T]$, where $\rho_{\text{YY}}[\mu, T]$ is the YY density for a uniform system. Similarly, the kinetic energy of the ground transverse mode is

$$E_{\text{kin},0} = \int dz \xi_{\text{kin}}[\mu(z), T],$$

where $\xi_{\text{kin}} = E_{\text{kin}}/\Delta z$ is the kinetic energy density of a uniform system of length $\Delta z$.

In the experiment [10] the system temperature was not sufficiently low that all transverse excitations were frozen out. We can account for this by observing that $g \rho_{0} \ll \hbar \omega_1$, so that the transverse excitations are well approximated as harmonic oscillator states with energies $\hbar \omega_1$ ($j = 0, 1, 2, ...$), where we have removed the zero-point energy $\hbar \omega_1$. The transversely excited states can then be accurately described as independent ideal 1D Bose gases with chemical potentials $\mu_j(z) = \mu(z) - j \hbar \omega_1$ [10]. Accounting for the degeneracy factor $j + 1$, the 1D position- and momentum-space densities for the atoms in transversely excited states are, respectively,

$$\rho(z) = \sum_{j=1}^{\infty} \frac{j+1}{\Lambda T} g_{1/2}(2) e^{\beta \mu_j(z)};$$

$$n_c(k_z) = \sum_{j=1}^{\infty} \frac{2}{\pi} \left\{ e^{\beta \frac{k_z^2}{2 \rho_{0}(z)} - \mu_j(z)} - 1 \right\}^{-1},$$

where $\beta = 1/k_B T$, and $\Lambda_T = (2 \pi \hbar^2 / m k_B T)^{1/2}$ is the thermal de Broglie wavelength. Using Eq. (4) we obtain the kinetic energy of the excited modes, $E_{\text{kin},c}$ [cf. Eq. (1)]. Combining these results, the full YY description of the trapped quasi-1D system is given by the total density, $\rho(z) = \rho_0(z) + \rho_c(z)$, total atom number $N = \int \rho(z) dz$, and the total average kinetic energy per particle $E_{\text{kin},t}/N = (E_{\text{kin},0} + E_{\text{kin},c})/N$.

Before comparing the YY thermodynamic predictions for $E_{\text{kin},t}/N$ with the experimental measurements of Ref. [10], we outline the second theoretical approach used here—the SPGPE—which allows for the determination of both global thermodynamic quantities and the full momentum distribution in the weakly interacting regime. In the SPGPE approach, the system field operator for the lowest transverse mode is split into two parts, $\hat{\Psi}_0(z) = \psi_C(z) + \psi_I(z)$, representing coherent ($\psi_C$) and incoherent ($\psi_I$) regions [17]. The coherent region is defined by an energy cutoff $\epsilon_{\text{cut}}$ such that it contains all highly occupied modes. It can then be described as a classical field (i.e. $\psi_C \rightarrow \psi_C$) evolving according to the simple evolution equation.

$$i \hbar \frac{\partial}{\partial t} \hat{\Psi}_0 = \hat{H}_0 \hat{\Psi}_0,$$

where $\hat{H}_0 = \int dz \left( -\frac{1}{2} \nabla^2 + V(z) + \frac{\mu}{k_B T} \right) \hat{\Psi}_0$. The Wigner function of the SPGPE can be written as

$$W(z, \rho, \phi) = \langle \hat{\Psi}_0(z) \hat{\Psi}_0^\dagger(z') \rangle,$$

where $\langle \cdot \rangle$ denotes the average over the initial state of the system. The density matrix of the SPGPE is given by

$$\rho_{\text{SPGPE}} = \int dz_1 dz_2 W(z_1, z_2, \rho, \phi) \hat{\Psi}_0(z_1) \hat{\Psi}_0^\dagger(z_2),$$

and the kinetic energy density is

$$E_{\text{kin}} = \int dz \left( \frac{1}{2} \rho(z) \beta \frac{\hbar^2}{\rho_{\text{kin}}(z)} - \mu(z) \right),$$

where $\rho_{\text{kin}}(z) = \int dz' W(z, z', \rho, \phi)$ is the kinetic density of the SPGPE.
where $L$ density can be sampled in both position and momentum. The values of $d$ distributions (black dots), bimodal Gaussian fits (red dashed lines) and theoretical SPGPE best-fits (blue solid lines) yielding the temperature $T$ and kinetic energy $E_{\text{kin},t}$. The bimodal Gaussian fits are difficult to distinguish from the SPGPE fits, but give slightly different values for the total kinetic energy (see inset of Fig. 3). The inset shows the atom number (open circles, left axis) and modal Gaussian fits (circles) [10], SPGPE fits (triangles) [29], YY thermometry (crosses), and the classical ideal gas model (gray line). The temperature is obtained from the width of the broad Gaussian component using a classical ideal-gas model. We note that there is very little quantitative difference between the heuristic bimodal Gaussian momentum fits and those based on the microscopic SPGPE formalism. However, there is a distinct difference in the temperatures extracted using the two methods (see Fig. 3).

In Ref. [10] the experimental momentum distributions were fitted with the sum of two Gaussians (see Fig. 2), giving the full momentum distribution and total kinetic energy. Examples of best-fit momentum distributions obtained using the SPGPE approach [27] are compared to the experimental data [28] in Fig. 2. We find quantitative agreement throughout the crossover from the nearly ideal Bose gas to the weakly interacting quasicondensate regime. From these fits we can determine both the kinetic energy per particle and the temperature for the data.

In Ref. [10] the experimental momentum distributions were fitted with the sum of two Gaussians (see Fig. 2), giving the total number of atoms $N$, and the total kinetic energy $E_{\text{kin},t}$. The temperature $T$ was obtained from the width of the broadest Gaussian component using a classical ideal-gas model. We note that there is very little quantitative difference between the heuristic bimodal Gaussian momentum fits and those based on the microscopic SPGPE formalism. However, there is a distinct difference in the temperatures extracted using the two methods (see Fig. 3).

We can also use the YY formalism for thermometry in this system, as there is a one-to-one correspondence between the temperature $T$ and the kinetic energy per particle $E_{\text{kin},t}/N$ for a given atom number $N$. We compare the temperature estimates of YY thermometry, the SPGPE momentum fits [29], and the broad Gaussian fits as a function of $E_{\text{kin},t}/N$ in the

\[ \rho_C(z) = \int_{E_C} \frac{dk_z}{2\pi} \frac{1}{e^{\beta \left[ \epsilon_{k_z} + 2g\rho_C(z)|\mu| \right]} - 1}. \]  

(6)

where $\epsilon_{k_z} = \hbar^2 k_z^2 / 2m + V(z)$ [23]. The incoherent region momentum density $n_I(k_z)$ is obtained by a similar procedure, and the transverse ground mode distribution is $n_{II}(k_z) = n_C(k_z) + n_I(k_z)$. Atoms in the excited transverse modes are treated as for the YY formalism, thus giving the full momentum distribution and total kinetic energy.

Figure 2. (Color online) Examples of the experimental momentum distributions (black dots), bimodal Gaussian fits (red dashed lines) and theoretical SPGPE best-fits (solid blue lines) yielding the temperature $T$ and kinetic energy $E_{\text{kin},t}$. The bimodal Gaussian fits are difficult to distinguish from the SPGPE fits, but give slightly different values for the total kinetic energy (see inset of Fig. 3). The inset shows the atom number (open circles, left axis) and modal Gaussian fits (circles) [10], SPGPE fits (triangles) [29], YY thermometry (crosses), and the classical ideal gas model (gray line). The temperature is obtained from the width of the broad Gaussian component using a classical ideal-gas model. We note that there is very little quantitative difference between the heuristic bimodal Gaussian momentum fits and those based on the microscopic SPGPE formalism. However, there is a distinct difference in the temperatures extracted using the two methods (see Fig. 3).

We can also use the YY formalism for thermometry in this system, as there is a one-to-one correspondence between the temperature $T$ and the kinetic energy per particle $E_{\text{kin},t}/N$ for a given atom number $N$. We compare the temperature estimates of YY thermometry, the SPGPE momentum fits [29], and the broad Gaussian fits as a function of $E_{\text{kin},t}/N$ in the

Figure 3. (Color online) Temperature of a quasi-1D Bose gas in a harmonic trap (with $\omega_x / 2\pi = 8.5$ Hz and $\omega_z / 2\pi = 3280$ Hz [23]) as a function of the total kinetic energy $E_{\text{kin},t}$ determined by bimodal Gaussian fits (open circles) [10], SPGPE fits (triangles) [29], YY thermometry (crosses), and the classical ideal gas model (gray line). The inset shows the atom number (open circles, left axis) and $E_{\text{kin},t}$ (right axis) as determined by bimodal Gaussian fits (open circles) [10] and SPGPE fits (triangles) to the momentum distributions as a function of the final rf evaporation frequency $\nu_{\text{rf}}$.
main panel of Fig. 5. A key result of this paper is that the YY temperatures collapse to the same curve as that obtained from the SPGPE fits to the full momentum distributions.

The inset of Fig. 3 shows the total number of atoms as a function of the final rf frequency for evaporative cooling, as well as a comparison of $E_{\text{kin},t}$ from the bimodal Gaussian fits and the SPGPE momentum fits. The temperature estimates from the SPGPE and YY methods differ for each final $\omega_{\text{rf}}$ due to small differences between the SPGPE and bimodal Gaussian fits (see the comparison in Fig. 5). However, if the YY thermometry utilised the values of $E_{\text{kin},t}$ determined from the SPGPE fits, the temperature estimates would be identical. Thus, the uncertainty in the temperature is due almost entirely to the experimental uncertainty in determining $E_{\text{kin},t}$. We note that the temperatures determined here typically lie above the estimates from the Gaussian fits of Ref. [10] — this illustrates the improved sensitivity of YY thermometry and SPGPE theory.

The utility of YY thermometry is that it relies on a single measurement—the kinetic energy per particle—which can be readily obtained via focusing and straightforward density imaging. This provides a simpler alternative to thermometry based on the measurement of density fluctuations of the gas [3]. It does not require any prior theoretical knowledge of the full momentum distribution, which is a challenging task in the strongly-correlated regime.

Finally, the results we present in this paper provide further quantitative validation of the SPGPE in the regime of current experiments with 1D quasi-condensates (see also [19]). As the SPGPE approach is suited to both equilibrium and dynamical simulations, this opens up an exciting avenue for exploring non-equilibrium phenomena in this system (e.g. quenches [30]), which cannot be explored using the YY solutions or equilibrium quantum Monte Carlo techniques.

In summary, we have studied the momentum properties of a finite-temperature 1D Bose gas, and given exact results for the rms width of the momentum distribution. We have outlined a procedure from which this simple quantity can be used for sensitive kinetic-energy thermometry of a quasi-1D harmonically trapped Bose gas using the YY formalism and the LDA. This method is applicable to all temperatures and repulsive interaction strengths. We have also performed a full characterization of the momentum distribution in the weakly interacting regime, and made a quantitative comparison with experimental data. Given the importance of momentum-space analysis in cold-atom research to date, our results provide a more complete picture of the homogeneous and trapped 1D Bose gas, and open up new prospects for measuring and probing these systems in the strongly-correlated regime.

MJD and KVK acknowledge support by the ARC through the Discovery Project scheme (DP110101047, DP1094025). PBB acknowledges support by Marsden contract 09-UOO-0924 and FRST contract NERF-UOOX0703. The work of AvA and NJvD is supported by FOM and NWO. NJvD acknowledges D. Stamper-Kurn for stimulating discussions.
[25] P. B. Blakie and M. J. Davis, Phys. Rev. A, 72, 063608 (2005).
[26] This is consistent with the conclusions of J.-B. Trebbia et al., Phys. Rev. Lett. 97, 250403 (2006), who found that the Hartree-Fock theory failed to describe the entire 1D Bose gas, i.e., the combination of both the coherent and incoherent regions.
[27] The best-fit SPGPE momentum distributions were determined by calculating the function $n(k_z, \mu, T)$, and performing a least-squares fit to the experimental data constrained to the experimentally determined atom number.
[28] The experimental momentum distributions were obtained by applying a focusing pulse to the trapped atoms, followed by a rapid switch off of the trapping potential [10]. The cloud then rapidly expands in the radial direction, effectively switching off the interactions, while it is still contracting axially. The axial contraction can therefore be treated as free propagation so that the measured spatial distribution in the focus is converted to a momentum distribution via the time-of-flight transformation $\hbar k_z = mz/\tau$. The timescale $\tau$ is calculated using the ABCD matrix formalism for matter-wave propagation giving $\tau = 15$ ms with an estimated systematic uncertainty of 10% [for a detailed discussion, see Ref. [24], in particular Section 4; see also J.-F. Riou et al., Phys. Rev. A 77, 033630 (2008)]. A very similar method was used to obtain the momentum distribution of a (nearly) 2D Bose gas: S. Tung, G. Lamporesi, D. Lobser, L. Xia, E. A. Cornell, Phys. Rev. Lett. 105, 230408 (2010).
[29] The two lowest $N$ data points in Fig. 3 use $\epsilon_{\text{cut}} = 150\hbar\omega_z$, and the remainder with $N < 17 \times 10^3$ use $\epsilon_{\text{cut}} = 300\hbar\omega_z$. For $N \gtrsim 17 \times 10^3$, $\mu$ is sufficiently negative that few modes are appreciably occupied and the coherent region vanishes. The system is then entirely described by Hartree-Fock theory.
[30] C. N. Weiler, T. W. Neely, D. R. Scherer, A. S. Bradley, M. J. Davis, and B. P. Anderson, Nature, 455, 948 (2008).