Abstract: The ionospheric equivalent slab thickness (τ) is defined as the ratio of the total electron content (TEC) to the F2-layer peak electron density (NmF2), and it is a significant parameter representative of the ionosphere. In this paper, a comprehensive statistical analysis of the diurnal, seasonal, solar, and magnetic activity variations in the τ at Guam (144.86°E, 13.62°N, 5.54°N dip lat), which is located near the magnetic equator, is presented using the GPS-TEC and ionosonde NmF2 data during the years 2012–2017. It is found that, for geomagnetically quiet days, the τ reaches its maximum value in the noontime, and the peak value in winter and at the equinox are larger than that in summer. Moreover, there is a post-sunset peak observed in the winter and equinox, and the τ during the post-midnight period is smallest in equinox. The mainly diurnal and seasonal variation of τ can be explained within the framework of relative variation of TEC and NmF2 during different seasonal local time. The dependence of τ on the solar activity shows positive correlation during the daytime, and the opposite situation applies for the nighttime. Specifically, the disturbance index (DI), which can visually assess the relationship between instantaneous τ values and the median, is introduced in the paper to quantitatively describe the overall pattern of the geomagnetic storm effect on the τ variation. The results show that the geomagnetic storm seems to have positive effect on the τ during most of the storm-time period at Guam. An example, on the 1 June 2013, is also presented to analyze the physical mechanism. During the positive storms, the penetration electric field, along with storm time equator-ward neutral wind, tends to increase upward drift and uplift F region, causing the large increase in TEC, accompanied by a relatively small increase in NmF2. On the other hand, an enhanced equatorward wind tends to push more plasma, at low latitudes, into the topside ionosphere in the equatorial region, resulting in the TEC not undergoing severe depletion, as with NmF2, during the negative storms. The results would complement the analysis of τ behavior during quiet and disturbed conditions at equatorial latitudes in East Asia.

Keywords: τ; geomagnetic equator; magnetic storm

1. Introduction

The ionospheric equivalent slab thickness (EST, usually named τ) is defined as the ratio of the total electron content (TEC) to the F2-layer peak electron density (NmF2) and it thus represents an imaginary equivalent depth of the ionosphere, with a constant uniform density of the F2 peak. Strictly speaking, when TEC is obtained by means of GNSS (Global Navigation Satellite System) signals, the τ value will contain a plasmaspheric component in addition to the ‘pure’ ionospheric component [1]. In the view of satellite-to-ground radio communications, τ is a useful parameter in that it includes information on both the topside and bottomside ionosphere. The τ is, thus, very helpful in understanding the nature of variations of the upper atmosphere and is, therefore, capable of addressing many
useful ionospheric parameters, such as Chapman scale height $H$, topside ‘half-peak density height’ $h_{0.5\text{top}}$, and temperature changes in the ionosphere/thermosphere systems [2–7]. In addition, according to the definition of $\tau$, it connects VTEC and $NmF2$, belonging to different fields of remote sensing, which can be used to deduce $NmF2$ from the TEC, and vice versa [8–13].

The ionospheric slab thickness has been a subject of study since the 1960s [8–33]. They found that the diurnal, seasonal, and solar activity variations of the $\tau$ show significant dependence on the location of the observing station. The greatest variability in $\tau$ is observed during periods of geomagnetic storms, and the influence of geomagnetic activity on the $\tau$ shows different patterns as the observing station and solar activity vary. Modelling studies have also been carried out to provide global $\tau$ [8–13,26,29].

The climatology of $\tau$ near the magnetic equator has been mentioned few times in previous studies. Odeyemi et al. [24,25] investigated the corresponding morphologies of ionospheric profile and peak parameters over an equatorial station in west Africa, Ilorin ($8.50^\circ$N, 4.68°E), during a low solar activity year 2010, and they found that the post-sunset increase in $\tau$ is more prominent during solstice than equinoxes, of which the December solstice perceives a relatively higher magnitude than the June solstice in 2010. Àlágbé [30] studied the correlation between the $\tau$ and $B0$ at the Burkina Faso station ($12.4^\circ$N, 1.5°W, dip $5.9^\circ$N), and they found that the maximum of $\tau$ at the station appeared in the daytime both high and low years of solar activity. The first peak occurred at between 05LT–06LT in the low solar activity year, and the daytime peak appears at about 11LT; the peak at sunrise in the high solar activity year is not apparent, while the daytime peak occurs at 12LT–13LT. The value of $\tau$ in high solar activity year is generally higher than that of low solar activity year. Duarte-Silva et al. [31] studied the $\tau$ at Palmas and São José’s dos Campos during one year of extremely low solar activity (from March 2009 to February 2010), and these two places are located at the inner edge of the anomaly region and the southern crest of the anomaly, respectively. They found that the $\tau$ at Palmas begins to increase gradually until its maximum value is reached between 13LT–15LT, and the maximum average peak values of $\tau$, of about 477 km, were observed over Palmas during daytime (08LT–16LT) throughout the December solstice months. In addition, the study of $\tau$ behavior during the storm time at equatorial latitudes is relatively less. Therefore, long-term monitoring and analyzing observations at equatorial latitudes are needed to give a more comprehensive study of the climatology of $\tau$ and geomagnetic storm effect on it at this latitude.

Guam station (144.86°E, 13.62°N, 5.54°N dip lat) is located at the equatorial north latitudes, where it lies between western Pacific Ocean and east of the Eurasian continent, and the surrounding area is dominated by marine environment. The ionospheric data are relatively sparse at this region, and there are few specific studies investigating the $\tau$ in this region. In order to strengthen the understanding of the variation of ionospheric thickness in East Asia near the magnetic equator, this paper uses the GPS-TEC and $foF2$ data at Guam station, between 2012 and 2017, to statistically analyze $\tau$ dependence with season, local time, and solar activity. Specifically, we investigate the effect of geomagnetic storms on the $\tau$ variation in this station. This paper is organized as follows: Section 2 briefly describes the data and methodology employed in this study. In Section 3, $\tau$ variation is presented, with an emphasis on its variations during geomagnetic storms. Section 4 discusses the derived results and use one example to analyze the causes of storm enhancement of $\tau$. Finally, conclusions of our study are found in Section 5.

2. Data and Methods of Analysis

The GPS TEC data in this paper are derived from UNAVCO (University NAVSTAR Consortium) database (http://www.unavco.org/ accessed on 23 November 2021). After downloading the receiver-independent exchange data (RINEX format) from UNAVCO, the GPS-TEC program developed by Seemala (https://Seemala.blogspot.com accessed on 23 November 2021) is applied to derive TEC values and elevation angle cutoff of 30° is adopted in order to eliminate the multi-path effects. In addition, the ionospheric shell...
height, used to retrieve TEC, is 350 km. Uwamahoro et al. [34] explained, in detail, how the GPS-TEC software works. Moreover, the GPS-TEC software used in this study has been compared with other techniques, such as the one presented in Ref. [35], alongside the European Geostationary Navigation Overlay System (EGNOS) algorithm, which was taken as a reference in Ref. [36]. Generally, both software is consistent with EGNOS algorithm, but the Gopi software was found to be closer to EGNOS in low-latitudes. It is worth noting that the same GPS-TEC software used in the current work has extensively been used to derive TEC in ionosphere and in τ climatology studies that involved TEC computation [25,37–43]. Moreover, we used the regional kriging interpolation method to convert the STEC obtained by software to VTEC, rather than directly use the average values of VTEC data calculated by the software.

Simultaneous foF2 data are obtained from the GIRO (Global Ionospheric Radio Observation) database (http://umler.uml.edu/DIDbase/ accessed on 23 November 2021) [44]. In the data processing, we filtered out the ionosonde data, which had a confidence score (CS) of less than 80 (maximum 100). The NmF2 values are computed from the critical frequency of the F2 layer (foF2) scaled from the recorded automatically ionograms, where $NmF2 = 1.24 \times 10^{10} (foF2)^2$. Then, the TEC and NmF2 data, with a 15 min resolution, are used to compute the equivalent τ by using the following relation:

$$\tau = \frac{TEC}{NmF2} = \frac{TEC}{1.24 \times 10^{10} (foF2)^2}$$ (1)

where TEC is measured in TEC units ($10^{16}$ electrons/m²) and foF2 is the critical frequency of the F2-layer, with foF2 given in MHz and τ in meters. For simplicity, we transform it in kilometers in the following.

Since (dayside) ionospheric density is produced, mainly, by solar EUV radiation (which is inferred by F10.7), we use F10.7 index to study the effect of solar activity on τ, and they are downloaded from NOAA (ftp://ftp.swpc.noaa.gov/pub/indices/old_indices/ accessed on 23 November 2021). Figure 1 shows the overall pattern of the solar F10.7 index during the period 2010–2017. It can be seen from Figure 1 that the F10.7 in 2014 was the largest among these 8 years, with an average value of 145.9 SFU (Solar Flux Unit). The F10.7 in 2010, 2016, and 2017 was small, with the average values being 80 SFU, 88.7 SFU, and 77.3 SFU, respectively. Due to the very high rate of data availability in 2016 (94.6%), this paper selects 2016 as the year of low solar activity and 2014 as the year of high solar activity to study the influence of solar activity on τ.

![Figure 1. F10.7 variations during the years 2012–2017.](image-url)
Previous studies have found that geomagnetic storms have effect on the $\tau$ variation \([15,27,32,33]\). In order to obtain the $\tau$ variation during geomagnetic quiet condition, we selected all the geomagnetic storms satisfying $\text{Dst}_{\text{min}} < -30$ during the period 2010–2017, according to the Dst (http://wdc.kugi.kyoto-u.ac.jp/ accessed on 23 November 2021) index, and the $\tau$ data, during the selected geomagnetic storm, are excluded from the overall $\tau$ dataset. Then, the $\tau$ data during the geomagnetic quiet condition are divided into 288 grids, according to January–December and 01–23LT. Recently, Pignalberi et al. \([22]\) used similar methodology to analyze $\tau$ at mid latitudes. The seasons are defined as $\pm 45$ days around vernal and autumnal equinoxes, June solstice, and December solstice.

The geomagnetic index Dst is used in this paper to select geomagnetic storm with minimum Dst < $-50$. To quantitatively describe the overall pattern of the geomagnetic storm effect on the $\tau$ variation, we define the disturbance index (DI):

$$\text{DI} = \frac{\tau_s}{\tau_m} - 1$$ (2)

where $\tau_s$ and $\tau_m$ are storm-time and monthly median $\tau$, respectively. Figure 2 shows an example during 2013 in Guam. Figure 2a illustrates the Dst variation, and the arrow ‘main phase onset (MPO)’ marks universal time of the magnetic main phase onset at 18:00 UT on 6 June. Figure 2b shows the variation of measured $\tau$ (blue lines) during 6–8 June and corresponding $\tau$ median (red) in June. Figure 2c presents the calculated DI index according to the measured $\tau$ value and its Median value. In addition, DI index in (storm time 1) S1 time period is used to study the characteristics of DI index during geomagnetic storm, while DI index in time periods (quiet time 1) Q1 and (quiet time 2) Q2 are used to study the one during geomagnetic quiet condition.

![Figure 2](image-url) An example for the calculation of DI index during the period 6–8 June 2013.

3. Results

According to the grid division of January–December and 0LT–23LT, Figure 3 gives the distribution of mean and standard deviation of $\tau$ on this grid. It can be seen from the figure that the first peak of $\tau$ appears at 11LT–12LT in all seasons, and it is the maximum value throughout the day. Moreover, there is a second peak that appears at 18LT–22LT in winter and equinox. During the nighttime, $\tau$ is significantly smaller than that in the daytime. From the perspective of seasonal variations, the $\tau$ in winter and equinox are generally greater than that in summer. In addition, the annual maximum appeared in
October 12LT, reaching 437.6km, and the annual minimum appeared in April 5LT, reaching 105.2km. Moreover, the \( \tau \) shows great variability during the night-time, especially during the pre-sunrise and post-sunset period, which is consistent with previous studies [7,22].

Figure 3. Map of mean (a) and standard deviation (b) of equivalent slab thickness (\( \tau \)) under geomagnetically quiet condition, according to January–December and 0LT–23LT division in 2012–2017.

Figure 4 compared the distribution of mean and standard deviation of \( \tau \) in the year of high solar activity and low solar activity, based on the same grid division. It can be seen from the figure that the EST in high/low activity year show some common features. That is, the first peak of \( \tau \) appears from 11LT to 12LT in all season. The \( \tau \) in 0LT–10LT is smaller than that in other periods, both during the years of high or low solar activity. The results also demonstrate that the \( \tau \) during the daytime in winter and equinox is larger than the \( \tau \) in summer. As for the differences between \( \tau \) during the high and low solar activity year, it is worth noting that the \( \tau \) in the 10LT–22LT time period of the high solar activity year seems greater than the average value of all years, while for other time periods, the \( \tau \) is smaller. The opposite situation applies for the change in low solar activity year. In addition, nighttime and solar terminator hours are also characterized by the highest dispersion, in the same way as Figure 3b. Specifically speaking, \( \tau \) shows more variability during the post-sunset period in the high solar activity year, while \( \tau \) shows more variability during the pre-sunrise period in the low solar activity year.

Among previous studies of \( \tau \) variation near geomagnetic equator, Álágbé. [30] found the maximum of \( \tau \) at Burkina Faso (12.4°N, 1.5°W, dip 5.9°N) appeared during the day, regardless of the solar activity of the year. The daytime peak appears at about 11LT in the low solar activity year; it appears at 12–13LT in the high solar activity year. Duarte-Silva et al. [31] found that the \( \tau \) at Palmas station (10.12°S, 48.21°W, 7.73°S dip) in the low solar activity have the following rules: the \( \tau \) begins to increase gradually after sunrise, and its maximum value appeared between 13LT to 15LT. Moreover, annual maximum of \( \tau \) reached 550 km in summer (that is, December) at the station. The \( \tau \) pattern in this paper is similar to these findings, but it is not completely consistent. The ionospheric longitudinal difference may contribute to the differences between them [45–47].

Figure 5 presents the diurnal variations of mean and standard deviation of \( \tau \) in summer, winter, and equinox at a geomagnetically quiet condition. It can be seen from Figure 5 that the first peak of \( \tau \) in winter is significantly higher than that in summer, and the equinox \( \tau \) is almost the same as the winter \( \tau \) during this period (range 2 (R2) interval). There is another peak of \( \tau \) in post sunset period in winter and equinox, and the summer \( \tau \) is also significantly smaller than that in winter and equinox during this period (range 3 (R3) interval). Moreover, the \( \tau \) during the post-midnight period is the smallest in equinox, especially during the period of 03LT–06LT (range 1 (R1) interval). In previous studies, the
sunrise peak of $\tau$ is a hot topic [21,26,33]. However, Figure 5 shows that, starting from 04LT in winter and summer, and 05LT in equinox, the $\tau$ gradually increases, and there is no peak. There is only a small peak appears at 06LT in winter. In addition, it can also be seen from the figure that the nighttime is characterized by the high dispersion, especially during the pre-sunrise and post-sunset period. During the pre-sunrise period, $\tau$ has highest variability in the summer, and $\tau$ has highest variability in the equinox during the post-sunset period.

![Figure 4](image_url)

**Figure 4.** Map of mean and standard deviation of equivalent slab thickness ($\tau$) under geomagnetically quiet condition. (a) Mean of $\tau$ in high solar activity year 2014; (b) Mean of $\tau$ in low solar activity year 2016; (c) Standard deviation of $\tau$ in high solar activity year 2014; (d) Standard deviation of $\tau$ in low solar activity year 2016.

Taking into account the rapid changes during the sunrise and sunset period, mean daytime (08LT–16LT) and night-time (20LT–04LT) values of $\tau$, for magnetically quiet days during the solar maximum phase 2014 and the solar minimum phase 2016, for three seasons (summer, winter, and equinox) are calculated (Table 1) to investigate the dependence of $\tau$ on solar activity. It can be seen from the table that, regardless of the high and low years of solar activity, the $\tau$ during the day is greater than the $\tau$ at night, and the seasonal characteristics of daytime $\tau$ (the $\tau$ in winter and equinox are greater than the summer $\tau$) are more obvious in the years of high solar activity. The most important finding is that the mean $\tau$ during the day, in the high solar activity, is greater than that in the low solar activity, whereas the opposite situation applies for the $\tau$ at night.
Figure 4. Map of mean and standard deviation of equivalent slab thickness (τ) under geomagnetically quiet condition. (a) Mean of τ in high solar activity year 2014; (b) Mean of τ in low solar activity year 2016; (c) Standard deviation of τ in high solar activity year 2014; (d) Standard deviation of τ in low solar activity year 2016.

Figure 5 presents the diurnal variations of mean and standard deviation of τ in summer, winter, and equinox at a geomagnetically quiet condition. It can be seen from Figure 5 that the first peak of τ in winter is significantly higher than that in summer, and the equinox τ is almost the same as the winter τ during this period (range 2 (R2) interval). There is another peak of τ in post sunset period in winter and equinox, and the summer τ is also significantly smaller than that in winter and equinox during this period (range 3 (R3) interval). Moreover, the τ during the post-midnight period is the smallest in equinox, especially during the period of 03LT–06LT (range 1 (R1) interval). In previous studies, the sunrise peak of τ is a hot topic [21,26,33]. However, Figure 5 shows that, starting from 04LT in winter and summer, and 05LT in equinox, the τ gradually increases, and there is no peak. There is only a small peak appears at 06LT in winter. In addition, it can also be seen from the figure that the nighttime is characterized by the high dispersion, especially during the pre-sunrise and post-sunset period. During the pre-sunrise period, τ has highest variability in the summer, and τ has highest variability in the equinox during the post-sunset period.

Table 1. Mean daytime (08:00–16:00LT) and night-time (20:00–04:00LT) values of ionospheric equivalent slab thickness (τ) for magnetically quiet days during the high solar activity year 2014 and the low solar activity year 2016.

| Year | Seasons | Mean Daytime Values of τ (km) | Mean Nighttime Values of τ (km) |
|------|---------|------------------------------|-------------------------------|
| 2014 | Winter  | 388.9                        | 232.1                         |
|      | Summer  | 303.9                        | 188.9                         |
|      | Equinox | 364.6                        | 250.7                         |
| 2016 | Winter  | 299.3                        | 255.5                         |
|      | Summer  | 288.6                        | 238.5                         |
|      | Equinox | 301.2                        | 268.4                         |

The ionosphere is not the same every day, even under undisturbed geomagnetic and solar conditions. In order to compare storm-time τ variation with quiet time τ variations. Figure 6 displays the mean DI index during quiet-time. In this case, the index is

\[ \text{DI} = \frac{\tau_q}{\tau_m} - 1 \]

where \( \tau_q \) and \( \tau_m \) are quiet-time and monthly median τ, respectively. According to if the DI is positive or negative, the daily variation of mean and standard deviation of DI index are shown in two panels, respectively. The blue curves represent the positive DI variation, and the red curves represent the negative DI variation. It is important to note that thick lines correspond to the average values of positive and negative DIs for the entire period, and vertical lines correspond to STD. As it shown, the τ is most stable in the daytime (08LT–19LT), and the DI index ranges from \(-0.12\) to \(0.14\) during this period. On the contrary, the τ is more variable during the period of night-time to sunrise period (22LT–07LT) with DI index more/less than \(0.19/\-0.15\). The τ is most variable during the sunrise period, with maximum/minimum having reached \(0.46/\-0.3\). It is consistent with the above results that the great variability of τ at sunrise, indicating that the τ values fluctuate sharply around sunrise. Previous observations and model simulations also illustrate that the day-to-day variability reaches the maximum at dawn, and they suggested that the day-to-day variability of neutral winds in the E region (≤~130 km) is the primary driver of the day-to-day variability of dawn ionosphere at the geomagnetic equator [48–51].
For the post sunset period (19LT–22LT), the DI increases/decreases from 0.12/−0.11 to
0.23/−0.15, whereas it decreases/increases from 0.21/−0.2 to 0.11/−0.11 during the post
sunrise period (07LT–08LT).

Figure 6. Diurnal variation of the mean and standard deviation (STD) of DI index at geomagnetically
quiet conditions.

Figure 7 presents the storm-time DI variation. The storm-time DI pattern is basically
consistent with that of quiet-time but with more variability. During the daytime
(08LT–19LT), the τ has smaller variability, and the DI indices are between −0.2 to 0.27. By
contrast, the τ is more variable during the nighttime to sunrise period (23LT–07LT), with
DI indices more/less than 0.287/0.168. In addition, the τ is also most variable during the
sunrise period (05LT–07LT), and maximum/minimum DI index are −0.45 and 0.88, respecti-
vively. For the post-sunset period (19LT–21LT), the DI increases/decreases from 0.22/−0.15
to 0.44/−0.17, and it decreases/increases from 0.26/−0.19 to 0.15/−0.12 during the sunrise
period (07LT–08LT).

Figure 7. Same as Figure 6 but for geomagnetic storm periods.
To better illustrate the geomagnetic effect on the $\tau$, Figure 8a compares the storm-time and quiet-time DI variations. One can find that the storm-time DI is more perturbed than quiet-time DI. Most of the storm-time positive DI are larger than quiet-time positive DI values, and this phenomenon is more apparent during night-time, especially during sunrise period. On the other hand, the storm-time negative DI are a little smaller than quiet-time negative DI values, except the pre-dawn period (4:30LT–6:30LT) with large difference. In order to illustrate the difference between them, we calculated the DI difference based on the sign of the DI index, and the result is shown in Figure 8b. As can be seen from Figure 8b, the positive DI difference during the storm is below 0.106 from 08LT to 19LT during the day. It is relatively large at sunrise (04LT–07LT) and sunset (19LT–21LT), reaching 0.45 and 0.24, respectively. The negative DI difference is less than $-0.08$ in all time periods, except for the period of sunrise, which reached $-0.15$ between 05LT–06LT. Therefore, we preliminarily draw the conclusion that geomagnetic storms have positive effects on the $\tau$ during most of the storm period, except they have negative effects on the $\tau$ during the sunrise period (4:30LT–6:30LT) in some cases.

Figure 8. (a) Comparison of diurnal DI index for geomagnetically quiet condition and geomagnetic storm condition (b) The minus between the DI index during geomagnetic storm periods and DI index at geomagnetically quiet time, according to the sign of DI.
4. Discussion

The Guam station is located near the magnetic equator. The most important factor affecting local ionospheric state is the fountain effect and the ionospheric equatorial anomaly (EIA) caused by it. At the same time, the geographic latitude of Guam Station is 13.43°N, and the trans-equator neutral wind can affect its ionospheric state significantly. Strong geomagnetic storms lead to positive or negative ionospheric storms, depending on season, local time, and the universal time (UT) arrival time of storms [52–55]. Different scales of neutral atmosphere waves, such as gravity waves, thermospheric tides, and planetary waves, driven by lower atmosphere processes, can also influence the ionosphere directly through plasma transport, or indirectly through electrodynamic processes [56,57]. Therefore, the τ in this paper displays a complicated pattern. Considering that geomagnetic activity is the main contributor to the relative ionospheric variability, we first discussed the τ under geomagnetic quiet conditions in Guam.

In order to explain the diurnal variation of τ, Figure 9 gives the quiet-time variation of TEC and NmF2 in Guam. In addition, it can be easily determined that the local time variation of τ depends on the variation in TEC and NmF2, obtained from Equation (1), as the following:

$$\frac{d\tau}{dt} = \frac{1}{NmF2} \frac{dTEC}{dt} - \frac{TEC}{NmF2^2} \frac{dNmF2}{dt}$$  \hspace{1cm} (3)

![Figure 9](image)

**Figure 9.** (a) Diurnal variations of mean and standard deviation of total electron content (TEC) in Guam during winter, summer, and equinox, respectively. (b) Diurnal variations of mean and standard deviation of F2-layer peak electron density (NmF2) in Guam during winter, summer, and equinox, respectively.

To better understand the possible formation mechanisms of the peak of τ, the temporal variation in the ionospheric parameters, TEC and NmF2, are given in Figure 10. According to Equation (3), the increase in the ratio variation of τ can be divided into following categories: (1) TEC increase and NmF2 decrease; (2) TEC and NmF2 increase, but the ratio of TEC increase is larger than that of NmF2; (3) TEC and NmF2 decreases but the ratio of NmF2 decrease is larger than that of TEC. The opposite situation applies for the decrease in τ. As shown in Figure 9, TEC and NmF2 consistently increase with the solar zenith angle, which increases after sunrise. Meanwhile, new electrons produced by the photoionization, in the magnetic equatorial ionosphere, began to be transported upward and diffused downward due to the fountain effect. Considering that the photoionization plays a much more important role in controlling the overall variability in the integrated quantity (i.e., TEC) than a localized quantity (i.e., NmF2) [58], the increases of ratio in TEC are larger than that in NmF2, as presented by Figure 10, and it causes the increases in τ during the forenoon.
Illustrated in Figure 9. As described above, the absence noontime bite-out in TEC is absent in summer, and equinox, respectively. Figure 10. (a) Variation rates in TEC for winter, summer, and equinox, respectively. (b) Variation rates in NmF2 for winter, summer, and equinox, respectively.

Figure 5 shows that the $\tau$ in Guam reaches its maximum value at noontime, consistent with previous results for equatorial latitudes [24,25,30,31]. Moreover, Figure 5 also shows that the first peak of $\tau$, in winter and equinox, are larger than in summer and this phenomenon is a new finding of this study. It is observed from Figure 9b that $NmF2$ in winter and equinox show ‘noon bite-out’ during the period of first $\tau$ peak, whereas TEC do not show this phenomenon during the full year (R2 interval). Apparently, these variations, together, result in seasonal minimum of the first peak in summer, according to Equation (1). Noon bite-out in the low latitude to low equator has long been known, and previous studies illustrate that this feature is mainly caused by the upward $E \times B$ drifts in the F region during daytime and noon bite-out was absent in the daytime TEC variation [59,60]. Lee [61] used TEC, $NmF2$, and $hmF2$ data, observed by the Jicamarca Digisonde and SAMI2, to model the noon bite-out under geomagnetic quiet condition and found that the reason why the noontime bite-out in TEC is absent is, principally, the upward $E \times B$ drifts because the drifts mainly account for the absent noontime bite-out in bottomside TEC and topside TEC. As for the absence of noon bite-out in summer, it might be caused by the neutral winds in Guam. Chen et al. [62] also suggested that the reason is seasonally dependent trans-equator plasma transport, induced by neutral winds. Since the meridional winds are equatorward during the summer daytime in Guam, plasma outflow, induced by trans-equator transport, results in significantly electron density depletion before the fountain effect induced bite-out. As a result, the pre-noon $NmF2$ peak is restrained. Thus, noontime bite-out variation pattern does not appear in summer in Guam.

As the solar zenith angel decreases and vertical drift velocity takes a downward trend, the electron density continues to decrease after noontime. This indicates a reduction in the $NmF2$ and TEC after midday, as illustrated in Figure 9. As described above, the photoionization plays a much more important role in controlling the overall variability in the TEC than $NmF2$. The variation of $NmF2$ is, therefore, less sensitive to the solar zenith angle compared to TEC. As shown in Figure 10, the ratio in decrease in TEC is larger than that in $NmF2$, which causes the decrease in $\tau$ in the afternoon.

Figure 5 depicts that there is another peak of $\tau$ that appears around sunset in winter and equinox. As shown in Figure 9, TEC and $NmF2$ continue to decrease during this period, due to the solar radiation decreases (R3 interval). It is the fact that the ratio of decreasing in TEC is smaller than that in $NmF2$ results in the increase in $\tau$, based on Equation (3) and Figure 10. It is generally accepted that pre-reversal enhancement (PRE), which refers to the strong upward $E \times B$ drift during the post-sunset period, appears from the equator to the low latitude region. Moreover, the electrodynamic process, such as $E \times B$, has stronger and more direct control only in the F-region height and density. Ionosphere generally exists at heights between 60–1000 km, whereas the plasmasphere could reach tens of thousands of kilometers. Therefore, the plasmasphere receives more solar radiation than ionosphere.
after sunset, and TEC thus decreases slower than \( NmF2 \). Based on these scenarios, there is an increase in \( \tau \) during sunset period. In addition, night-time enhancement in electron density, at low latitudes, has been studied by several groups of workers and the post-sunset enhancement is known to be caused by the PRE. It is necessary to note the night-time enhancement in electron density is said to occur if its value is greater than the exponentially decaying background value. Su et al. [63] reported that the frequency of occurrence of the post-sunset TEC enhancement has two maxima in equinox, and these equinoctial maxima are separated by a broad summer minimum at the eastern low-latitude station Taiwan (11.5° geomagnetic latitude). Therefore, the post-sunset \( \tau \) does not illustrate an increase in summer, whereas they show another peak in winter and equinox.

Agrawal et al. [64] has shown that pre-reversal increase in the vertical \( E \times B \) drift, the primary source for the electron density enhancement, is strongest at equinox in both the eastern and western longitudes, and the night-time enhancement in \( NmF2 \) and TEC is, therefore, greatest in equinox, as displayed by Figure 9. However, there is no solar radiation after sunset, and \( F \) region only constitutes a part of space environment when evaluating TEC. Therefore, the TEC enhancement is not apparent as \( NmF2 \) enhancement. It causes the equinox/solstice ratio in \( NmF2 \) to grow larger than that in TEC during the post-midnight to sunrise period, as shown in Figure 9. Figure 9 also shows that the post-midnight enhancement is most obvious in summer, consistent with previous studies. As the recombination process takes place steadily, one can find that the TEC and \( NmF2 \) give a more pronounced contribution than \( F2 \) during the pre-sunrise period. Therefore, the \( \tau \) in equinox is smallest, which is caused by the aforementioned largest equinoctial \( NmF2 \) and the faster decrease in TEC in equinox than in solstice, as shown in Figures 9 and 10, respectively.

In previous studies, the sunrise peak of \( \tau \) has been a hot topic [21,22,26,33]. These studies indicate that the pre-sunrise peak is a widely observed feature from low to high latitudes, and \( \tau \) even reaches a maximum at sunrise for specific seasons. However, Figure 5 shows that, starting from 04LT in winter and summer, and 05LT in equinox, the \( \tau \) gradually increases, and \( \tau \) gets very low values (around 100 km) during the pre-sunrise period. For equatorial latitudes, the similar phenomenon was also seen for other longitudes, as follows. Añávé [30] studied the correlation between the \( \tau \) (TF2 in his paper) and B0 at the Burkina Faso station (12.4°N, 1.5°W, dip 5.9°N), and they found that the \( \tau \) at Burkina Faso station (12.4°N, 1.5°W, dip 5.9°N) seemed to also have a very low value (less than 100 km) at pre-sunrise period. Duarte-Silva et al. [31] have also found that the Palmas station (10.12°S, 48.21°W, 7.73°S dip lat), located in the inner edge of the anomaly region, also seems to get very low values around the pre-sunrise period. Odeyemi et al. [24,25] have studied the \( \tau \) over Africa Ilorin (8.50°N, 4.68°E), and they found that the \( \tau \) also seems to get very low values on some months during the low solar activity period. Therefore, the very low values at pre-sunrise hours might not be a special phenomenon in Guam because it could also be seen at other longitudinal equatorial latitude station. Moreover, the TEC should give a more pronounced contribution than \( NmF2 \) during the pre-sunrise period. Therefore, the \( \tau \) begins to increase during the pre-sunrise period, as can be seen in Figure 5, but it does not reach a peak. To further understand the morphology of \( \tau \) during the pre-sunrise period, more observations and modelling are needed to provide a statistical picture and study the physical mechanism of this phenomenon in the future.

Among previous studies on \( \tau \), the dependence of \( \tau \) on solar activity has shown different correlations in different latitudes [16,19]. Añávé [30] shows that the day-time \( \tau \) on the magnetic equator, in the high solar activity, is significantly greater than that in the low solar activity. Jayachandran et al. [21] also found that \( \tau \) at noon is positively correlated with solar activity when studying \( \tau \) in low and middle latitudes. The results shown in Table 1 confirm their conclusions. As for the \( \tau \) decreases with solar activity at night, it should be due to the fact the H+O+ transition height in the low solar activity is lower than
that in the high solar activity at night. It is known that the $\tau$ is sensitive to the variations of $\text{H}/\text{O}$ ratio at the F2 peak or equivalent to the transition level at which $[\text{O}] = [\text{H}]$. Large downward fluxes of $\text{H}$ can decrease the $\text{O}$ to $\text{H}$ transition levels, thereby increasing the topside content and, hence, the slab thickness.

Geomagnetic activity is another key factor affecting the perturbations of the $\tau$. Previous studies have shown that the influence of geomagnetic activity on the $\tau$ depends on the location of the observing station and solar activity. For the mid latitude, the $\tau$ seems to systematically enhance during periods of geomagnetic disturbances [15,32,65]. However, magnetic activity does not appear to have significant influence on the $\tau$ variations at low latitude [22,66]. Until now, there is no specific statistical study analyzing the geomagnetic storms effect on the $\tau$ at the near equator latitudes (best to our knowledge). Therefore, we made a statistical study of the geomagnetic storms effect on the $\tau$ in Guam, based on the aforementioned DI index during the geomagnetic storm period from the year 2010 to 2017.

Comparing with the geomagnetic quiet time DI index shown in Figure 6, the results in Figure 8 suggest that the geomagnetic storm seems to have a positive effect on the $\tau$ during most of the storm period near the equatorial station Guam.

During geomagnetic storms, the main process affecting $\tau$ ($\text{TEC/}N_{m}F2$) in Guam includes: (1) penetration electric field driven by the solar wind ranging from high latitude to equator; (2) the equator-ward neutral wind resulting from particle precipitation and Joule heating at high latitude, sometimes accompanied with travelling atmosphere disturbance (TAD); (3) the disturbance dynamo electric fields produced by the globally altered thermospheric winds during magnetic storms; (4) composition changes, driven by storm time, altered neutral winds. These coupled drivers together with background thermosphere-ionosphere generate ionospheric storms and control their strengths in complicated ways. In the following, we will start from these aspects and use an example on the 1 June 2013 in Guam to analyze why the $\tau$ tends to increase when the positive or negative ionospheric storm occurs.

Figure 11 shows the geophysical conditions during 1–2 June 2013. As IMF Bz turns southward at 0:00UT, Dst began to decrease sharply, reaching a minimum of $-124$ nT at 8:00UT, and then, the geomagnetic storm entered the recovery phase. The DI indices of $\tau$, TEC and $N_{m}F2$ of 1 June are given in Figure 12. As shown in the figure, with the start of the main phase, caused by southward reversal of IMF Bz, a TEC positive storm occurred during the main phase, but the increase in $N_{m}F2$ during this period was very small, resulting in a $\tau$ positive storm during this period (T1 interval in Figure 12).
Figure 12. (a) The DI index of TEC and \(NmF2\) on 1 June 2013. (b) The DI index of \(\tau\) on 1 June 2013.

It is generally accepted that the penetration electric field and equatorward neutral wind should be the primary drivers of the positive disturbance at this latitude \([43,58]\). From the fact that TEC positive storm occurred during the main phase of the magnetic storm, caused by the southward reversal of IMF BZ in Figure 12, we infer that there is a penetrating electric field during this period. The daytime penetrating electric field tends to move the plasma from the equatorial region into higher latitudes through the fountain effect. Meanwhile, the increased upward vertical drifts also transport local plasma in the bottomside ionosphere or the F2 layer to higher altitudes where the chemical recombination rate is low, so plasma accumulates in the topside ionosphere. On the other hand, the plasma, transported upward from the bottomside ionosphere, would be compensated by daytime ionization. As a result, a large increase in TEC is accompanied by a rather small increase in \(foF2\) and \(\tau\) tends to enhance during the positive ionospheric storm, as shown in T1 interval. In order to confirm above theory, Figure 13 shows the DI index of ionospheric bottomside TEC (BTEC) and topside TEC (TTEC) relative to their monthly median value on 1 June. Among them, BTEC is obtained by integrating the electron density below the ionospheric peak height \(\text{hm}\), while TTEC is obtained by subtracting BTEC from the TEC. Figure 13 shows that the ionospheric positive storm in this time period is mainly caused by the TTEC, while the BTEC gives little contribution. From the perspective of plasma scale height, which has a close relationship with \(\tau\), the increased upward drifts changed the shape of the topside ionosphere and increased the topside effective plasma scale height \(dh/d(ln(\text{Ne}))\) significantly, so the whole \(\tau\) is thus increased. Moreover, the enhanced equatorward wind during storm time should also play an important role in producing changes in electron densities, as it can raise the F2 region to a higher altitude to inhibit field-aligned ambipolar diffusion, thus causing changes in the shape of the ionospheric density profile \([67,68]\).

As the geomagnetic storm entered the recovery phase, \(NmF2\) experienced a negative storm, but the disturbance of TEC during this period was small, resulting in a positive disturbance of \(\tau\) (T2 interval in Figure 12). It is well known that the disturbance dynamo electric fields produced by the globally altered thermospheric winds during magnetic storms and composition changes brought by equatorward neutral winds are the main drivers of negative storms in the region we are interested in \([54]\). However, TEC and \(NmF2\) might have contrasting behavior during the recovery phase of the geomagnetic storm, as shown in Figure 12. Since the topside ionosphere, at equatorial latitudes, connected with middle and low latitudes of the region near the F2 peak at the flux tube, the storm time enhanced equatorward wind tends to push more plasma at low latitudes from the region near the F2 peak into the topside ionosphere in the equatorial region \([69]\). Consequently, the topside TEC in the equatorial region can undergo an obvious enhancement, and it
causes \( \tau \) to increase, combining with the negative disturbance of \( NmF2 \). Figure 13 confirms this theory, BTEC experienced a large negative disturbance, but TTEC did not, and even was in a positive disturbance state for most of the time during T2 interval. Therefore, the positive disturbance of TTEC, caused by equatorial neutral wind during a storm, makes TEC decrease little when \( NmF2 \) has a negative storm occur, resulting in an increase in \( \tau \).

![Figure 13. The DI index of TTEC and BTEC on 1 June 2013.](image)

5. Conclusions

Based on the TEC and \( NmF2 \) data from the years 2012–2017, this paper statistically analyzed the \( \tau \) at equatorial latitude Guam. The results show great diurnal, seasonal, solar, and geomagnetic activity variation. A brief review of observations made by other researchers has also been presented, and we obtained the following results, which confirm similar studies:

1. The peak of \( \tau \) appeared at noon, consistent with previous studies on equatorial latitudes.
2. There is a post-sunset peak in \( \tau \) observed during the winter and equinox, and it means that \( NmF2 \) is decreasing faster than TEC, which can be associated with the higher post-sunset TEC enhancement occurrence in equinox/winter, proving previous conclusions. In addition, the \( \tau \) continues to decrease after post-midnight and the equinoctial one is smallest due to the largest nighttime \( NmF2 \) enhancement in equinox and the faster decrease in TEC in equinox than that in solstice.
3. The dependence of \( \tau \) on the solar activity are different for daytime and nighttime: the daytime \( \tau \) seems to increase with solar activity, as TEC is more sensitive to the solar activity than \( NmF2 \), whereas the nighttime one decreases with solar activity at night, and it should be due to the fact the H+O+ transition height in the low solar activity is lower than that in the high solar activity at night.
4. The \( \tau \) has more variability during nighttime than daytime, during both the geomagnetically quiet and disturbed conditions, and the greatest variability of \( \tau \) appeared at sunrise.

Moreover, we obtained some new results, which provide interesting insights into \( \tau \) of this region:

1. The \( \tau \) at noon is larger in winter and equinox than in summer. It is probably due to the absence of \( NmF2 \) noontime bite-out in summer at this region.
2. There is no pre sunrise peak in \( \tau \) and \( \tau \) get low values during the pre-sunrise period at Guam. However, previous studies indicate that pre sunrise peak is a widely observed feature, from low to high latitude, and \( \tau \) even reaches maximum at sunrise for specific seasons. The contradiction is probably due to Guam being located at equatorial latitude, as the low values in the pre sunrise period could also be seen at...
other longitudinal equatorial latitude station. In addition, longitudinal difference might also contribute to the difference.

3. The geomagnetic storm seems to have a positive effect on the $\tau$ during most of the storm period in Guam, except at sunrise period, when the $\tau$ attains large variability, even at the geomagnetically quiet condition. This study also provides a new physical explanation for the observed effect of geomagnetic storm on $\tau$ in Guam. During the positive storms, the penetration electric field along with storm time equator-ward neutral wind tends to increase upward drift and uplift F region, causing the large increase in TEC accompanied by relatively small increase in $NmF2$. On the other hand, an enhanced equatorward wind tends to push more plasma, at low latitudes, into the topside ionosphere in the equatorial region, resulting in the TEC, which does not undergo severe depletion as $NmF2$ does during the negative storms. Therefore, the geomagnetic storm seems to enhance $\tau$ both during the positive and negative ionospheric storms.

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