Scintillation Properties of Ti- and Zr-doped Gd$_2$Hf$_2$O$_7$ Crystals

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Zr-doped and Ti-doped Gd$_2$Hf$_2$O$_7$ single crystals were synthesized by the optical floating zone (FZ) technique to evaluate their scintillation properties. The synthesized single crystals show scintillation signals with a broad peak at 510 nm, and scintillation decay curves consist of two exponential functions with decay time constants on the nanosecond order. Both the Ti and Zr doping processes seem to enhance the scintillation properties of Gd$_2$Hf$_2$O$_7$. Regarding their afterglow characteristics, all the Gd$_2$Hf$_2$O$_7$ crystals show similar afterglow levels after X-ray irradiation for 2 ms.

1. Introduction

Scintillators are phosphors that immediately convert absorbed high-energy photons$^{(1)}$ or particles$^{(2,3)}$ into numerous UV–visible photons, and they have been utilized in many fields of radiation detection such as medical imaging$^{(4)}$ security,$^{(5)}$ astrophysics,$^{(6)}$ and geophysical resource exploration.$^{(7)}$ For X- or γ-ray detection, scintillators generally require high scintillation light yield (LY), short decay time, high energy resolution, large effective atomic number, high density, and low afterglow levels; however, since no developed scintillators satisfy all these requirements, users must select a suitable one from existing scintillators. Thus, many researchers have made effort to search and develop new desirable materials.

In particular, a large effective atomic number and a high density provide a high interaction probability between materials and ionizing radiations,$^{(8)}$ so pyrochlore rare-earth hafnate is one of the potential choices as scintillators with high density (7.93–9.95 g/cm$^3$)$^{(9)}$ and large effective atomic number ($Z_{\text{eff}} = 64–69$).$^{(10)}$ For scintillator applications, some pyrochlore rare-earth hafnates have been investigated only in a polycrystalline form including ceramics and powders.$^{(11–15)}$ In general, scintillators have been developed in the single crystal form owing to the high optical transparency and homogeneity; however, the growing need is the driving force of developing new scintillators. In the case of rare-earth hafnate, since crystal growth is quite difficult because of the high melting point (~2400 °C), few studies on the growth of pyrochlore rare-earth hafnate compounds have been reported.$^{(16–18)}$ Previously, we introduced the optical floating zone (FZ) method as an effective and powerful technique for searching new
oxide scintillators with high melting points.\textsuperscript{(19–23)} Up to now, pyrochlore $RE_2Hf_2O_7$ ($RE = $ La, Gd, Lu) single crystals have been prepared using an optical FZ furnace equipped with Xe lamps, and the scintillation properties of these single crystals were reported for the first time by Nakauchi \textit{et al.}\textsuperscript{(18)}

To improve scintillation characteristics by increasing oxygen vacancies, Zr and Ti ions with radii smaller than that of the Hf ion are introduced into $RE_2Hf_2O_7$ single crystals to distort the crystal lattice because Hf-based oxide materials typically show luminescence due to oxygen vacancies such as F and F$^+$ centers.\textsuperscript{(24,25)} Such investigations have been performed for several scintillator candidates showing intrinsic luminescence, and activation has been achieved in intrinsic scintillators.\textsuperscript{(26)} Enhancements by the substitution of homologous ions were also observed in $La_2Hf_2O_7$, and Zr and Ti ions improved the scintillation LY.\textsuperscript{(27)} In this paper, we focus on such an enhancement effect in pyrochlore Gd$_2$Hf$_2$O$_7$.

2. Materials and Methods

Before crystal growth, undoped, 1% Zr-doped, and 1% Ti-doped Gd$_2$Hf$_2$O$_7$ polycrystalline rods were synthesized by conventional solid-state reaction. First, raw powders of ZrO$_2$ (99.99\%, Furuuchi Chemical), TiO$_2$ (99.99\%, High Purity Chemicals), Gd$_2$O$_3$ (99.99\%, Furuuchi Chemical), and HfO$_2$ (99.95\%, Furuuchi Chemical) were mixed. Then, the homogeneously mixed powder was formed into cylinder rods by applying hydrostatic pressure. Furthermore, the obtained rods were heated at 1500 °C for 8 h in air. After that, crystal growth was performed by the optical FZ technique (FZ-T-12000-X-VPO-PC-YH, Crystal Systems).\textsuperscript{(18)} Here, the pulling-down and rotation rates were 30 mm/h and 3 rpm, respectively. To clarify the crystalline phase of the samples, the powder X-ray diffraction (XRD) pattern was measured using a diffractometer (MiniFlex600, Rigaku) over the $2\theta$ range from 5 to 90°.

To determine the scintillation properties of the samples, scintillation spectra under X-ray irradiation were measured using our original setup.\textsuperscript{(28)} The X-ray generator (XRB80N100/CB, Spellman) was equipped with an X-ray tube having a W anode target and a Be window. The X-ray tube was operated at a bias voltage of 80 kV and a tube current of 1.2 mA. In these analyses, the emitted scintillation photons from the samples were led to a spectrometer unit equipped with a CCD (DU-420-BU2, Andor) and a monochromator (Shamrock163, Andor) through an optical fiber. Here, the detector was cooled to 188 K using a Peltier module to reduce the thermal noise. Scintillation decay and afterglow curves were measured using an afterglow characterization system,\textsuperscript{(29)} and the photomultiplier tube (PMT) covered the spectral range from 160 to 650 nm. In these analyses, the decay time constant was calculated by least-squares fitting with the sum of two exponential decay functions. To measure pulse height spectra, a sample covered with four layers of Teflon reflectors was mounted on the window of the PMT (R7600-200, Hamamatsu Photonics). The scintillation signals were converted to electrical signals, which were processed by the preamplifier (113, ORTEC), shaping amplifier (570, ORTEC), and multichannel analyzer (Pocket MCA 8000A, Amptek). Here, the shaping times for the synthesized samples and reference GS20 sample\textsuperscript{(30)} were 2 and 0.5 μs, respectively.
3. Results and Discussion

All the Gd₂Hf₂O₇ single crystals were successfully obtained, and the typical sizes of the as-prepared samples were 5–6 mm ⊙ × 15 mm. Figure 1 shows the powder XRD patterns of the Gd₂Hf₂O₇ single crystals. Compared with the reference pattern, only a single phase of Gd₂Hf₂O₇ with a cubic structure and a space group of Fd̅3m was observed. As illustrated in Fig. 2, the samples were cut into small pieces with sizes of 5–6 mm ⊙ × 3–5 mm thickness, and the top and bottom surfaces were polished to measure the luminescence properties.

Figure 3 shows the X-ray-induced scintillation spectra, and all the crystals exhibit scintillation signals with a broad peak at 510 nm. Although the emission origins have not been reported, we consider that the emission originated from oxygen vacancies related to an interpenetrating network of GdO₈ and HfO₆, as suggested in the case of La₂Hf₂O₇. The X-ray-induced scintillation decay curves are illustrated in Fig. 4, and the obtained curves are approximated by the sum of two exponential functions having decay time constants of 50–60 ns (~80%) and 600–900 ns (~20%). The values are similar to those reported for a polycrystalline form. The decay of the primary component is faster than that of a Bi₄Ge₃O₁₂ scintillator.
Figure 5 shows the pulse height spectra of the Gd$_2$Hf$_2$O$_7$ samples measured under $^{241}$Am α-ray (5.4 MeV), and full energy peaks can be observed in all the samples. In comparison with the channel of the observed full energy peak of the GS20 scintillator showing an LY of 4000 photons/MeV, (30) the scintillation LYs of the undoped, Zr-doped, and Ti-doped Gd$_2$Hf$_2$O$_7$ samples are respectively 140, 280, and 210 photons/MeV with a typical error of ±10%. Since both Zr- and Ti-doped samples show higher scintillation LYs than the undoped one, the Zr and Ti doping processes appear to enhance the scintillation output.

To prove the emission origin, X-ray-induced scintillation spectra were measured using the samples annealed under vacuum and in air, as illustrated in Fig. 6. The sample annealed in air shows a small peak shift to a wavelength longer than that in the case of the as-prepared sample. Although it is difficult to control oxygen defects by such a simple annealing process, this result suggests that at least the obtained emissions include those derived from oxygen defects such as F and F$^-$ centers as well as the other Hf-based oxides. (24,25)

The afterglow curves of samples with 2 ms X-ray irradiation are shown in Fig. 7. The afterglow levels are calculated from $100 \times (I_2 - I_{BG})/(I_1 - I_{BG})$, where $I_{BG}$, $I_1$, and $I_2$ are defined as the background signal intensity, the signal intensity during irradiation, and the signal

### Table 1

| Samples     | $\tau_1$ (ns) | $\tau_2$ (ns) |
|-------------|----------------|---------------|
| Undoped     | 54 ± 9 (81%)   | 817 ± 186 (19%) |
| Zr-doped    | 53 ± 8 (83%)   | 662 ± 129 (15%) |
| Ti-doped    | 50 ± 9 (86%)   | 603 ± 190 (16%) |

Figure 5. (Color online) Pulse height spectra measured using Gd$_2$Hf$_2$O$_7$ under $^{241}$Am α-ray (5.4 MeV) irradiation.

Fig. 6. (Color online) Scintillation spectra of undoped Gd$_2$Hf$_2$O$_7$ before (black) and after annealing under vacuum (blue) and in air (red).

Fig. 7. (Color online) Afterglow curves of Gd$_2$Hf$_2$O$_7$ single crystals after 2 ms X-ray irradiation.
intensity at $t = 20$ ms, respectively. The obtained afterglow levels of the undoped, Zr-doped, and Ti-doped crystals are 246, 346, and 236 ppm, respectively. All the samples exhibit similar values, but the Ti-doped one exhibits the best afterglow level. The values are close to that of Tl:CsI (~300 ppm) but higher than those of commercial oxide scintillators such as Ce:Gd$_2$SiO$_5$ (5.9 ppm).

4. Conclusions

In this study, we focused on pyrochlore Gd$_2$Hf$_2$O$_7$ as a scintillator candidate with high density. Undoped, Zr-doped, and Ti-doped Gd$_2$Hf$_2$O$_7$ single crystals were successfully synthesized by the optical FZ technique. All the samples exhibit scintillation with a broad emission band at 550 nm, and the results suggest that the origin of the emission is oxygen vacancies, because annealing in air affects the spectral shape. The doped samples show higher scintillation LYs than the undoped one, indicating that both Zr and Ti improve the scintillation properties of Gd$_2$Hf$_2$O$_7$.

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