Measuring single-top-associated Higgs production at the HL-LHC

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Abstract

Precision measurements of top-associated Higgs production are an important ingredient to unravel the \(CP\) nature of the Higgs boson. In this work, we constrain the \(CP\) nature of the top-Yukawa coupling taking into account all relevant inclusive and differential Higgs boson measurements. Based upon this fit, we show that it is crucial to disentangle single- and di-top-associated Higgs production for tightening indirect constraints on a \(CP\)-odd top-Yukawa coupling in the future. In this context, we propose an analysis strategy for measuring \(tH\) production at the HL-LHC without relying on assumptions about the Higgs \(CP\) character.

1 Introduction

After the discovery of a new particle which is consistent with the predictions for the Standard Model (SM) Higgs boson, the investigation of this particle is one of the main tasks for future LHC and High-Luminosity LHC (HL-LHC) runs. Especially in light of the absence of any direct evidence for beyond the SM (BSM) physics, the precise determination of the Higgs boson’s properties can give crucial hints to unravel the nature of BSM physics.

The determination of the Higgs boson’s \(CP\) properties is an important part of this program with a close connection to cosmology, since the Higgs boson could provide a new source of \(CP\) violation allowing to explain the baryon asymmetry of the Universe (BAU). While experimental studies have already excluded the possibility of the Higgs boson being a pure \(CP\)-odd state, the possibility of the Higgs boson being a \(CP\)-admixed state is so far less constrained.

The focus of the present study [1] is the by magnitude largest Higgs–fermion–fermion interaction: the top-Yukawa interaction. While the \(CP\) properties of the top-Yukawa coupling can also be constrained by electric dipole measurements or by demanding a sufficient amount of \(CP\) violation to explain the BAU (see e.g. [2]), we concentrate on the constraints imposed by current and future LHC measurements.
At colliders, $CP$-violating couplings can be constrained directly by measuring $CP$-odd observable. Measuring a non-zero value for such an observable would directly imply the presence of $CP$ violation. While proposals for $CP$-odd observables targeting the top-Yukawa coupling exists, their measurement is experimentally challenging. $CP$ violation in the Higgs–top-quark interaction, however, also induces deviations from the SM in $CP$-even observables. While a deviation from the SM in a $CP$-even observable is not guaranteed to be caused by $CP$ violation, a systematic investigation of indirect constraints is still a powerful method to narrow down the available parameter space for $CP$ violation.

In the present work, we perform a fit to derive bounds on a $CP$-violating top-Yukawa coupling following the indirect approach. We take into account all available inclusive and differential Higgs boson measurements. Based upon the results of this fit, we point out that a measurement of single top quark associated Higgs production — without relying on an assumption about the Higgs $CP$ character — would significantly enhance the sensitivity to a $CP$-violating top-Yukawa coupling. We then propose a strategy for performing such a measurements at the HL-LHC focusing on the Higgs boson decay to two photons.

2 Top-associated Higgs production as a probe of the top-Yukawa interaction

We parameterize BSM effects in the Higgs-boson interaction with top quarks in the form

$$L_{yuk} = -y_{SM}^{t} \sqrt{2} \tilde{t} (c_t + i \gamma_5 \tilde{c}_t) tH,$$

where $y_{SM}^{t}$ is the SM top-Yukawa coupling, $c_t$ rescales the $CP$-even top-Yukawa coupling ($c_t = 1$ in the SM), and $\tilde{c}_t$ constitutes a $CP$-odd top-Yukawa coupling ($\tilde{c}_t = 0$ in the SM). Additionally, we introduce the parameter $c_V$ rescaling the Higgs couplings to massive vector bosons. To parameterize the effect of additional BSM particles affecting the $H \rightarrow \gamma\gamma$ and $gg \rightarrow H$ processes, we, moreover, float the Higgs–gluon–gluon and the Higgs–photon–photon couplings as free parameters (denoted as $\kappa_g$ and $\kappa_\gamma$, respectively).

The most relevant processes to constrain the modified top-Yukawa coupling of Eq. (1) are Higgs production via gluon fusion, the Higgs decay to two photons, Z-boson associated Higgs production, and top-associated Higgs production. Top-associated Higgs production is of special interest since the top-Yukawa coupling appears at the tree level allowing for a comparably model-independent probe of the Higgs–top-quark interaction. Three different sub-processes contribute to top-associated Higgs production: $tH$, $tWH$, and $tH$ production. While $tWH$ and $tH$ production are negligible in the SM, their cross-section can be significantly enhanced in the presence of a $CP$-violating contribution to the top-Yukawa coupling. We do not consider constraints arising from processes involving a virtual Higgs boson (e.g. $t\bar{t}$ or $t\bar{t}t\bar{t}$ production).

We perform the global fit of this model to all relevant Higgs boson measurements — including the latest Higgs rate measurements from ATLAS and CMS as well as the $p_T$-binned simplified template cross-section measurements for the process $pp \rightarrow ZH, H \rightarrow b\bar{b}$ — using HiggsSignal [3–5].

In Fig. 1, we show two exemplary results of the described fit focusing on top-associated Higgs production. Since the different top-associated Higgs production channels are hard to distinguish experimentally, often a combination of these is measured. In the left plot of Fig. 1, the cross section for the sum of all three channels (normalized to the SM prediction) is shown as a function of $\tilde{c}_t$ as predicted by the fit described above. The dependence of $\mu_{tH+tWH+tWH}$ on $\tilde{c}_t$ is approximately flat indicating that even a future more precise measurement of $\mu_{tH+tWH}$ will not allow to tighten the constraints on $\tilde{c}_t$. If it would instead be possible to disentangle
Figure 1: Left: Signal strength for combined top-associated Higgs production for the five-dimensional global fit in dependence of $\tilde{c}_t$. The color encodes the profiled $\Delta \chi^2$ distribution of the fit. The white, light-gray, and dark-gray contour lines represent the 1\sigma, 2\sigma, 3\sigma confidence regions, respectively. Right: Same as left plot but the signal strength for $tH$ Higgs production divided by the signal strength for $t\bar{t}H$ and $tWH$ production is shown. The green band indicates the impact of the projected $tH$ measurement with 3 ab$^{-1}$ at the HL-LHC.

$tH$ production from $t\bar{t}H$ and $tWH$ production, the bounds on $\tilde{c}_t$ could be improved. This is evident in the right plot of Fig. 1 displaying the SM-normalized $tH$ over $t\bar{t}H + tWH$ cross section ratio as a function of $\tilde{c}_t$.

3 Measuring single-top-associated Higgs production

As discussed in Sec. 2, disentangling $tH$ and $t\bar{t}H + tWH$ production is an important step to tighten constraints on a $CP$-violating Higgs–top-quark interaction. One possible strategy is to exploit the different lepton multiplicities of $tH$ and $t\bar{t}H + tWH$ production [1] by defining a one-lepton and a two-lepton category. Since only $t\bar{t}H + tWH$ production can contribute to the two-lepton category, the cross section in the two-lepton category can be used as a $t\bar{t}H + tWH$ control measurement for the one-lepton category to which all three sub-channels contribute.

To enhance the sensitivity of this analysis, the one-lepton category is optimized for a high $tH$ event fraction. At this point it is important to make sure that such a $tH$ rate measurement is applicable for constraining the $CP$ nature of the Higgs boson. I.e., it should be ensured that the measurement is independent of the Higgs $CP$ nature (because otherwise the measurement could not be used to constrain the Higgs $CP$ nature).

The difficulty of this is indicated in Fig. 2 for the $H \rightarrow \gamma\gamma$ decay channel. While e.g. the rapidity difference between the leading $b$ jet and the leading non-$b$ jet $|\Delta y^{b-j}|$ is a good discriminator between $tH$ and $t\bar{t}H + tWH$ production (see upper left panel), cuts based on this observable introduce a large dependence of the measurement on the $CP$ character of the top-Yukawa coupling of $\sim 40\%$ (see upper right panel). Using instead the geometric mean of the leading jet and the Higgs boson rapidities $y^{j} \oplus y^{TT}$, $tH$ and $t\bar{t}H + tWH$ production can be disentangled equally well (see lower left panel) with a negligible ($\lesssim 2\%$) dependence on the Higgs $CP$ character (see lower right panel).

This strategy has been used to obtain a projection for a $tH$ measurement at the HL-LHC [1]: the $tH$ signal strength is projected to be below 2.21 at 95\% CL assuming SM-like data and using 3 ab$^{-1}$ (a limit five times stronger than the current strongest limit [6]). The constraint on $\tilde{c}_t$ which would be imposed by such a measurement (for SM-like data) is indicated by the green
4 Conclusion

In the present article, we investigated how a $CP$-violating Higgs–top-quark interaction can be constrained indirectly at the LHC. Using all relevant inclusive and differential Higgs-boson measurements, we performed a global fit showing that a more precise determination of top-associated Higgs production combining the $tH$, $t\bar{t}H$, and $tWH$ channels would not result in improved bounds on a $CP$-odd top-Yukawa coupling. In order to improve the sensitivity, it will be crucial to disentangle $tH$ from $t\bar{t}H + tWH$ production without relying on assumptions on the Higgs $CP$ character.

We then discussed a potential strategy for reaching this goal by defining signal regions with different lepton multiplicities. We showed that it is especially important for the optimization of the analysis cuts to always keep the dependence of the kinematic acceptances on the Higgs $CP$ character as low as possible. Using 3 ab$^{-1}$, we projected an upper limit on the $tH$ signal strength five times stronger than the current strongest limit.

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