Entropy generation analysis of peristaltic flow of magneto-nanoparticles suspended in water under second-order slip conditions

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Abstract. Peristalsis of magneto-nanoparticles suspended in water was investigated in this study. Fe\textsubscript{3}O\textsubscript{4}-water nanofluid was utilized for two-dimensional flow in channel, which was considered symmetric with compliant walls. Uniform magnetic field was applied and temperature equation was arranged for viscous dissipation. Also, second-order velocity and thermal slip conditions were utilized. Small Grashof number led to the perturbation solution. Examination of entropy generation was also carried out. Maxwell and Hamilton-Crosser models was used. The analysis was based on the comparative study of these two models representing the cylindrically and spherically shaped particles. Graphs for velocity, temperature, entropy generation, and Bejan numbers were plotted under the influence of sundry variables. Trapping was observed via plotting streamlines.

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1. Introduction

Nanotechnology has attracted the attention of recent researchers due to its ample applications to industrial, biomedical, and engineering fields. It has made many advancements in our daily life. For instance, through the use of nanoparticles, we can remove the harmful and dangerous viruses and bacteria from water. Vehicle fuel efficiency and corrosion resistance can also be improved by making the vehicle parts from nanocomposite materials. These materials can also be employed to produce sports goods with features such as higher strength, light weight, etc. The term nanofluid was first used by Choi [1]. Nanomaterials are very useful materials with particle sizes in the nano (10\textsuperscript{-9} meter) scale. Nanofluids are made by the mixture of a base fluid, e.g., water, oil, ethylene glycol, etc., and nanoparticles. Nanoparticles have different shapes, e.g., spherical, cylindrical, tube, blade, bricks, etc., and various types including the oxides, carbides, nanotubes, etc. Different two-phase models have been utilized by researchers in the literature. In two-phase models for the nanofluids, thermo-physical characteristics of the nanomaterial and base liquid are separately defined [2,3]. These models provide the advantage of investigating various types of nanofluids by using different nanoparticles and base fluids, hence the capability to make a comparison among them for identifying the ones with better performances. On the other hand, in single-phase models, like Buongiorno model [4], thermo-physical properties of each fluid are

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not specified. Thus, one can analyze the fluid in general by considering the prominent effects of nanofluid, i.e., Brownian motion and thermophoresis effects. Favorable thermo-physical properties of nanofluid have drawn the curiosity of researchers and have led to many advancements in engineering, biomedical, and industrial fields, hence being a hot topic of research. Mahanthesh et al. [5] investigated a two-phase model for nanofluids by utilizing the SWCNT and MWCNT in which the flow was made by a rotating disk. Hsiao [6] examined the stretching sheet problem for micropolar nanofluid. Makinde et al. [7] discussed the flow made by a rotating disk in the presence of aluminum and titanium alloy nanoparticles. Some other studies in the literature dealing with nanofluids can be found in [8–13], which shed more light on the utility of nanofluids.

In physiology, bulk fluid movement is due to the peristaltic phenomenon. In this mechanism, a progressive wave moves within the channel walls, forcing the fluid to flow. This phenomenon is significantly work in human body. Some obvious instances are the transportation of food from esophagus, movement of chyme, transportation of urine to bladder, transportation of lymph in lymphatic vessels, etc. In addition to physiological processes, peristalsis are significant in industrial processes and biomedical applications. Modern devices like dialysis machine, open heart bypass machine, etc. operate under peristalsis. This principle is useful for transferring corrosive fluids since it prevents their contamination with machinery. Hence, peristalsis is applicable to sanitary fluid transportation. Many engineering devices like hose pump, roller pump, finger pump, etc. operate on the same principle. The broad physiological, biomedical, and industrial applicability of peristalsis make them a significantly important area of research. Thus, many research works on peristalsis have been carried out after the pioneering research of Latham [14], e.g., [15–27]. Peristaltic transport of nanofluids can be consulted by many attempts in [28–36].

It is well known that several aspects such as chemical reaction, Joule heating, etc., in a thermodynamic system affect its entropy. Basically, the entropy of a system is a measure of disorder in it. Bejan [37] presented the first analysis of entropy. The existing literature provides insufficient information on peristaltic transportation of fluid with entropy (see [38–40]). Hence, the main objective of the present study is to further investigate this regime. Peristalsis of nanofluid subject to entropy generation under mixed convection is considered. Fe3O4 with water as the base fluid is used as the nanofluid. There is no study to deal with the second-order slip conditions in terms of both velocity and temperature and this study is aimed to fill this gap. Specifically, the Hamilton-Crosser and Maxwell models are comparatively used for the cylindrical and spherical particles, respectively. The resulting problems employing lubrication approach are solved by regular perturbation method. Grashof number is adopted as the perturbation quantity. Velocity, temperature, entropy generation, and Bejan number are sketched and examined. Trapping phenomenon is studied by graphical illustration of streamlines.

2. Flow configuration

Peristaltic flow of an incompressible nanofluid composed of Fe3O4 and water is considered. The channel (with the width of 2d) is considered symmetric. Flexible channel walls are placed at the positions y = ±η, where the left and right walls are denoted by - and +, respectively (see Figure 1). The rectangular coordinates system is set such that the x-axis lies in the direction of channel length and the y-axis is perpendicular to it. Temperature of the walls is maintained at T0. Contribution of the applied constant magnetic field is taken into account. Induced magnetic and electric fields effects are omitted. Mixed convection and viscous dissipation are studied. Sinusoidal wave has the wavelength of λ, amplitude of a, and speed of c. The shape of the wave is defined by the equation given below:

\[
y = \pm \eta(x, t) = \pm \left[ d + a \sin \frac{2\pi}{\lambda}(x - ct) \right].
\]  

(1)

The equations for the considered flow configuration are:

\[
\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0,
\]  

(2)

\[
\rho_{\text{eff}} \left( \frac{\partial}{\partial t} + u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y} \right) u = -\frac{\partial p}{\partial x} + \mu_{\text{eff}} \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) - \sigma_{\text{eff}} B_0^2 u + g(\rho \beta)_{\text{eff}} (T - T_0).
\]  

(3)

Figure 1. Flow geometry.
\[
\rho_{\text{eff}} \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right) = -\rho \frac{\partial p}{\partial y} - \sigma_{\text{eff}} B^2 v + \mu_{\text{eff}} \left[ \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right].
\]

(4)

\[
(\rho C)_{\text{eff}} \left( \frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} \right) = \mu_{\text{eff}} \left[ \left( \frac{\partial \rho}{\partial y} + \frac{\partial \rho}{\partial x} \right)^2 \right] \]

\[+2\left( \frac{\partial u}{\partial x} \right)^2 + \frac{\partial v}{\partial y} \right)^2 \right) + K_{\text{eff}} \left[ \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right].
\]

(5)

The quantities assumed in the above-mentioned equations are defined as: \(u\) and \(v\) for components of velocity in the \(x\) and \(y\) directions, respectively; \(\rho_{\text{eff}}\) effective density; \(p\) pressure; \(\mu_{\text{eff}}\) effective viscosity; \(g\) acceleration due to gravity; \(\sigma_{\text{eff}}\) effective thermal conductivity; and \((\rho C)_{\text{eff}}, K_{\text{eff}},\) and \((\rho C)_{\text{eff}}\) effective thermal expansion, effective thermal conductivity of nanofluids, and effective heat capacity, respectively. In addition, \(T\) stands for temperature and \(t\) represents time.

The notations \(\rho_{\text{eff}}, (\rho\beta)_{\text{eff}}, (\rho C)_{\text{eff}}, \mu_{\text{eff}},\) and \(K_{\text{eff}}\) for the two-phase model are:

\[
(\rho\beta)_{\text{eff}} = (1 - \phi)\rho_f \beta_f + \phi \rho_p \beta_p,
\]

\[
\rho_{\text{eff}} = (1 - \phi) \rho_f + \phi \rho_p,
\]

\[
(\rho C)_{\text{eff}} = (1 - \phi)(\rho C)_f + \phi(\rho C)_p,
\]

\[
\mu_{\text{eff}} = \frac{\mu_f}{(1 - \phi) \gamma_f^2}.
\]

\[
\sigma_{\text{eff}} = 1 + \frac{3 \left( \frac{\nu}{\nu_f} - 1 \right) \phi}{\left( \frac{\nu}{\nu_f} + 2 \right) - \left( \frac{\nu}{\nu_f} - 1 \right) \phi}.
\]

For Maxwell’s model:

\[
K_{\text{eff}} = \frac{K_p + 2K_f - 2\phi(K_f - K_p)}{K_p + 2K_f + \phi(K_f - K_p)}.
\]

For Hamilton-Crosser’s model:

\[
K_{\text{eff}} = \frac{K_p + (n - 1)K_f - (n - 1)\phi(K_f - K_p)}{K_p + (n - 1)K_f + \phi(K_f - K_p)},
\]

where the subscripts \(f\) and \(p\) represent the fluid and nanoparticles, and \(\phi\) depicts volume fraction of nanoparticles. Two models of effective thermal conductivity are used in the above equation. The Hamilton-Crosser model is employed for the cylindrical particles with \(n = 6\) and Maxwell model is adopted for spherical particles. Here, \(n\) represents the shape of the nanoparticles. It is defined by \(3/\Psi\), where \(\Psi\) depicts sphericity of nanomaterials. The amount of \(\Psi = 0.5\) is used for cylindrical-shape material and \(\Psi = 1\) for spherical-shape material.

Thermophysical properties of the base liquid and nanoparticles are mentioned in Table 1. Dimensionless parameters are:

\[
x^* = \frac{x}{\lambda}, \quad u^* = \frac{u}{c}, \quad y^* = \frac{y}{d},
\]

\[
\eta^* = \frac{\eta}{d}, \quad v^* = \frac{v}{c}, \quad t^* = \frac{ct}{\lambda},
\]

\[
p^* = \frac{\rho \nu^2}{\mu_f}, \quad \nu = \frac{T - T_0}{T_0}, \quad \text{Pr} = \frac{\mu_f C_f}{K_f},
\]

\[
\text{Re} = \frac{\rho \nu^2 d}{\mu_f}, \quad \text{Ec} = \frac{e^2}{C_f T_0}, \quad \text{Br} = \text{Pr Ec},
\]

\[
M = \sqrt{\sigma_f B_d d}, \quad \text{Gr} = \frac{g \rho \nu^2 T_0 d^2}{\mu_f},
\]

\[
u = \frac{\partial \psi}{\partial y}, \quad v = -\delta \frac{\partial \psi}{\partial x}.
\]

(7)

Here, \(\text{Pr}, \text{Re}, \text{Br}, \text{Ec}, M,\) and \(\text{Gr}\) denote the Prandtl, Reynolds, Brinkman, Eckert, Hartman, and Grashof numbers, respectively.

After applying the lubrication approach, we obtained:

\[
\frac{\partial p}{\partial x} = \frac{1}{(1 - \phi)^2} \frac{\partial^2 \psi}{\partial y^2} + \text{Gr} \nu \delta - M^2 A_1 \frac{\partial \psi}{\partial y},
\]

(8)

\[
\frac{\partial p}{\partial y} = 0.
\]

(9)

\[
K_f \frac{\partial^2 \theta}{\partial y^2} + \frac{\text{Br}}{(1 - \phi)^2} \left( \frac{\partial \psi}{\partial y} \right)^2 = 0,
\]

(10)

\[
A_1 = 1 + \frac{3 \left( \frac{\nu}{\nu_f} - 1 \right) \phi}{\left( \frac{\nu}{\nu_f} + 2 \right) - \left( \frac{\nu}{\nu_f} - 1 \right) \phi}.
\]

Table 1. Thermophysical parameters of water and nanoparticles [13]:

|   | \(\rho\) (kg m\(^{-3}\)) | \(C_p\) (J kg\(^{-1}\)K\(^{-1}\)) | \(K\) (W m\(^{-1}\)K\(^{-1}\)) | \(\beta\) (l/k) \times 10\(^{-6}\) | \(\sigma\) (\Omega m\(^{-1}\)) |
|---|----------------|----------------|----------------|----------------|----------------|
| Fe\(_3\)O\(_4\)| 5200 | 670 | 80.6 | 13 | 25000 |
| H\(_2\)O | 997.1 | 4179 | 0.613 | 210 | 0.05 |
\[ A_3 = 1 - \phi + \phi \left( \frac{(\rho\beta)_p}{(\rho\beta)_f} \right). \]

For Maxwell’s model:
\[ K_1 = \frac{K_p + 2K_f - 2\phi(K_f - K_p)}{K_p + 2K_f + \phi(K_f - K_p)}. \]

For Hamilton-Crosser’s model:
\[ K_1 = \frac{K_p + (n - 1)K_f - (n - 1)\phi(K_f - K_p)}{K_p + (n - 1)K_f + \phi(K_f - K_p)}. \]

Boundary conditions are:
\[ \frac{\partial \psi}{\partial y} \pm \beta_1 \frac{\partial^2 \psi}{\partial y^2} \pm \beta_2 \frac{\partial^3 \psi}{\partial y^3} = 0, \]
\[ \theta \pm \gamma_1 \frac{\partial \theta}{\partial y} \pm \gamma_2 \frac{\partial^2 \theta}{\partial y^2} = 0, \quad \text{at} \quad y = \pm \eta. \]

The corresponding stream function and temperature solutions are:
\[ \psi_0 = \frac{A_0 e^{\frac{-1}{2K_1}\sqrt{\frac{\pi}{M_1}}}}{A_1 M^2} + C_3 + yC_4, \]
\[ \theta_0 = -\frac{1}{4A_1 K_1 M^2} A_0 \left( e^{\frac{-1}{2\sqrt{\pi\eta}} - \frac{\pi}{2\eta}} \right) \left( C_2^2 + C_3^2 e^{\frac{-1}{\sqrt{\pi\eta}}} \right) + F_1 + yF_2. \]

3. Solution methodology

Perturbation technique is employed to achieve the solution. The small Grashof number is chosen as the perturbation parameter and the corresponding systems and their solutions are:

**Zeroth-order systems.** Here we have:
\[ \frac{1}{(1 - \phi)^{2\alpha}} \frac{\partial^2 \psi_0}{\partial y^2} - M^2 A_1 \frac{\partial^2 \psi_0}{\partial y^2} = 0, \]
\[ K_1 \frac{\partial^2 \theta_0}{\partial y^2} + \frac{Br}{(1 - \phi)^{2\alpha}} \left( \frac{\partial^2 \psi_0}{\partial y^2} \right)^2 = 0, \]
\[ \frac{\partial \psi_0}{\partial y} \pm \beta_1 \frac{\partial \psi_0}{\partial y^2} \pm \beta_2 \frac{\partial \psi_0}{\partial y^3} = 0, \quad \text{at} \quad y = \pm \eta. \]
\[ \left[ E_1 \frac{\partial^3}{\partial x^3} + E_2 \frac{\partial^3}{\partial x \partial y^2} + E_3 \frac{\partial^2}{\partial x \partial y} \right] \eta = \frac{1}{(1 - \phi)^{2\alpha}} \frac{\partial \psi_0}{\partial y^2} - M^2 A_1 \frac{\partial \psi_0}{\partial y}, \quad \text{at} \quad y = \pm \eta. \]
\[ \frac{\partial \psi_0}{\partial y} \pm \gamma_1 \frac{\partial \psi_0}{\partial y^2} \pm \gamma_2 \frac{\partial \psi_0}{\partial y^3} = 0, \quad \text{at} \quad y = \pm \eta. \]

2.1. Entropy analysis

Viscous dissipation is represented by:
\[ \Phi = \mu \varepsilon \left[ \left( \frac{\partial u}{\partial y} \right)^2 + \left( \frac{\partial v}{\partial x} \right)^2 \right]. \]

Volumetric entropy generation in dimensionless form is:
\[ S_{\text{gen}}^v = \frac{K_{\text{eff}}}{T_m} \left( \frac{\partial T}{\partial x} \right)^2 + \left( \frac{\partial T}{\partial y} \right)^2 + \frac{\Phi}{T_m}. \]

Dimensionless form becomes:
\[ N_s = \frac{S_{\text{gen}}^v}{S_G^v} = K_{\text{eff}} \left( \frac{\partial \theta}{\partial y} \right)^2 + \frac{Br}{\Lambda (1 - \phi)^{2\alpha}} \left( \frac{\partial \psi}{\partial y} \right)^2, \]
\[ S_G^v = \frac{K_{\text{eff}} T_0^2}{T_m^2 \beta}, \quad \Lambda = \frac{T_0}{T_m}. \]

Bejan number is:
\[ \text{Be} = \frac{N_{s\text{cond}}}{N_{s\text{cond}} + N_{s\text{visc}}}. \]

Eq. (15) comprises two parts; \( N_{s\text{cond}} \) is entropy generation due to temperature difference and \( N_{s\text{visc}} \) is entropy generation due to viscous dissipation.
\[ \psi_1 = -\frac{1}{24A_1^2 K_1 M^3} \left( A_0^4 \frac{A_2}{8} A_3 \text{Br} \frac{e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu}}{\sqrt{A_1}} \left( C_2^2 - C_2^2 e^{\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} \right) \right) \]

\[ - 12A_1 A_3 F_2 K_1 M^3 y^2 + 8A_0 A_1 M^3 \left( A_3 \text{Br} C_1 C_2 y^2 \right) \]

\[ - 3e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} K_1 \left( e^{\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} B_1 + B_2 \right) \]

\[ + y B_4 + B_3. \]  

(31)

\[ \theta_1 = \frac{1}{9A_1^2 K_1 M^3} A_0 \text{Br} \left( A_0^4 \frac{A_3}{2} \text{Br} C_1^2 e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} \frac{3 \sqrt{A_1} \text{M}}{3 \sqrt{A_1} \text{M}} \right) \]

\[ - \frac{A_0^4 A_3 \text{Br} C_1^2 e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu}}{3 \sqrt{A_1} \text{M}} K_1 M^2 \]

\[ - \frac{9}{2} A_0 A_1 B_2 C_2 e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} K_1 M^2 \]

\[ - \frac{9}{2} A_0 A_1 B_1 C_1 e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} K_1 M^2 \]

\[ - 9A_1^2 \left( B_2 C_1 + B_1 C_2 \right) K_1 M^4 y^2 \]

\[ - 3 \sqrt{A_0} A_3 C_2 e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} \left( \sqrt{A_0} \left( - 25A_0^3 \text{Br} C_1 C_2 + 6 \sqrt{A_1} F_2 K_1 M \right) \right) \]

\[ - 12A_0^2 \text{Br} C_1 C_2 y \]  

\[ - 3 \sqrt{A_0} A_3 C_1 e^{-\sqrt{\frac{\text{Br}}{A_1}} \text{M} \mu} \left( \sqrt{A_0} \left( 25A_0^3 \text{Br} C_1 C_2 + 6 \sqrt{A_1} F_2 K_1 M \right) \right) \]

\[ - 12A_0^2 \text{Br} C_1 C_2 y + G_1 + y G_2. \]

(32)

The constants Bi’s, Ci’s, Gi’s, and Fi’s can be computed via Mathematica.

4. Discussion

This section is devoted to the analysis of velocity, temperature, entropy generation, Bejan number, and stream lines. Each quantity is analyzed in a specific subsection.

4.1. Analysis of velocity

In this subsection, behavior of velocity is discussed under the influence of different important parameters. Figure 2 represents the impact of nanoparticle volume fraction on velocity profile. This graph shows decreasing behavior, which is due to the fact that by increasing the quantity of nanoparticles (\( \phi = 0.01, 0.03, 0.05, 0.07 \)), resistance to the fluid increases so that fluid velocity diminishes. In this study, the values for the Hamilton-Crosser’s model are greater than those for the Maxwell’s model. Figure 3 presents velocity against Hartman number. It demonstrates that velocity has decreasing behavior with increase in Hartman number \((M = 2, 3, 4, 5)\), because the Lorentz

![Figure 2](image-url)  

Figure 2. \( u \) in terms of \( \phi \) when \( t = 0.1, E_3 = 0.01, x = 0.2, E_2 = 0.01, E_1 = 0.02, M = 1.0, \varepsilon = 0.2, \)
Gr = 0.03, Br = 3.0, \( \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \) and \( \gamma_2 = -0.01. \)

![Figure 3](image-url)  

Figure 3. \( u \) in terms of \( M \) when \( t = 0.1, E_3 = 0.01, x = 0.2, E_2 = 0.01, E_1 = 0.02, \varepsilon = 0.2, \phi = 0.01, Br = 3.0, \)
Gr = 0.03, \( \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \) and \( \gamma_2 = -0.01. \)
force acts as a resistive force. Figure 4 deals with Grashof number, showing increase in velocity profile by enhancing Grashof number (Gr = 0.1–0.7). This is due to increase in buoyancy forces, which facilitate the flow. Velocity profile with wall parameters can be observed in Figure 5. The results illustrate that velocity increases with elastance parameters (E₁ = 0.01, 0.02) and (E₂ = 0.02, 0.04) and decreases with the damping parameter (E₃ = 0.01, 0.02). Obviously, elastance parameters cause lower resistance, leading to increase in velocity, whereas damping resists the flow. The results for the slip parameters are demonstrated in Figures 6 and 7. As shown, velocity profile shows enhancement when the slip parameters (β₁ = 0.1, 0.3, 0.5, 0.7) and (β₂ = −0.1, −0.3, −0.5, −0.7) are increased. It is also noteworthy that this behavior is more prominent with the second-order slip parameter than with the first-order one. Furthermore, the velocity profile is higher in Hamilton-Crosser’s model than in Maxwell’s model.

4.2. Analysis of temperature
In this subsection, temperature profile for different per-

Figure 4. $u$ in terms of $Gr$ when $t = 0.1$, $E₃ = 0.01$, $x = 0.2$, $E₂ = 0.01$, $E₁ = 0.02$, $ε = 0.2$, $M = 1.0$, $φ = 0.1$, $Br = 3.0$, $β₁ = 0.01$, $β₂ = −0.01$, $γ₁ = 0.01$, and $γ₂ = −0.01$.

Figure 5. $u$ in terms of $E₁$, $E₂$, $E₃$ when $t = 0.1$, $x = 0.2$, $ε = 0.2$, $M = 1.0$, $φ = 0.1$, $Br = 3.0$, $Gr = 0.03$, $β₁ = 0.01$, $β₂ = −0.01$, $γ₁ = 0.01$, and $γ₂ = −0.01$.

Figure 6. $u$ in terms of $β₁$ when $t = 0.1$, $E₃ = 0.01$, $x = 0.2$, $E₂ = 0.01$, $E₁ = 0.02$, $ε = 0.2$, $M = 1.0$, $φ = 0.1$, $Br = 3.0$, $Gr = 0.03$, $β₂ = −0.01$, $γ₁ = 0.01$, and $γ₂ = −0.01$.

Figure 7. $u$ in terms of $β₂$ when $t = 0.1$, $E₃ = 0.01$, $x = 0.2$, $E₂ = 0.01$, $E₁ = 0.02$, $ε = 0.2$, $M = 1.0$, $φ = 0.1$, $Br = 3.0$, $Gr = 0.03$, $β₁ = 0.01$, $γ₁ = 0.01$, and $γ₂ = −0.01$.

Figure 8. $θ$ in terms of $φ$ when $t = 0.1$, $E₃ = 0.01$, $x = 0.2$, $E₂ = 0.01$, $E₁ = 0.02$, $ε = 0.2$, $M = 1.0$, $Gr = 0.03$, $Br = 3.0$, $β₁ = 0.01$, $β₂ = −0.01$, $γ₁ = 0.01$, and $γ₂ = −0.01$. contend parameters is accounted for. Figure 8 provides graphs for $φ = 0.01, 0.03, 0.05, and 0.07$ versus temperature distribution. This graph illustrates that temperature profile diminishes with increasing $φ$. On the other hand, increase in $φ$ enhances thermal conductivity and
Figure 9. $\theta$ in terms of $M$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $\phi = 0.1$, $Br = 3.0$, $Gr = 0.03$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 10. $\theta$ in terms of $Gr$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $M = 1.0$, $\phi = 0.1$, $Br = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 11. $\theta$ in terms of $E_1$, $E_2$, $E_3$ when $t = 0.1$, $x = 0.2$, $\varepsilon = 0.2$, $M = 1.0$, $\phi = 0.1$, $Br = 3.0$, $Gr = 0.03$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 12. $\theta$ in terms of $\gamma_1$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $M = 1.0$, $\phi = 0.1$, $Br = 3.0$, $Gr = 0.03$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, and $\gamma_2 = -0.01$.

Figure 13. $\theta$ in terms of $\gamma_2$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $M = 1.0$, $\phi = 0.1$, $Br = 3.0$, $Gr = 0.03$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, and $\gamma_1 = 0.01$.

that temperature remains higher for spherical-shape particles than for cylindrical-shape ones.

4.3. Analysis of entropy generation and Bejan number
This subsection accounts for entropy generation and Bejan number for different embedded parameters. Fig-
Figure 14. $N_s$ in terms of $\phi$ when $t = 0.1, E_3 = 0.01, x = 0.2, E_2 = 0.01, E_1 = 0.02, \varepsilon = 0.2, M = 1.0, BrA^{-1} = 1.0, Gr = 0.03, Br = 3.0, \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \text{and } \gamma_2 = -0.01.$

Figure 15. $N_s$ in terms of $M$ when $t = 0.1, E_3 = 0.01, x = 0.2, E_2 = 0.01, E_1 = 0.02, \varepsilon = 0.2, \phi = 0.1, BrA^{-1} = 1.0, Gr = 0.03, Br = 3.0, \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \text{and } \gamma_2 = -0.01.$

Figure 16. $N_s$ in terms of $Gr$ when $t = 0.1, E_3 = 0.01, x = 0.2, E_2 = 0.01, E_1 = 0.02, \varepsilon = 0.2, \phi = 0.1, M = 1.0, BrA^{-1} = 1.0, Br = 3.0, \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \text{and } \gamma_2 = -0.01.$

Figure 17. $N_s$ in terms of $BrA^{-1}$ when $t = 0.1, E_3 = 0.01, x = 0.2, E_2 = 0.01, E_1 = 0.02, \varepsilon = 0.2, \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \text{and } \gamma_2 = -0.01.$

Figure 18. $N_s$ in terms of $E_1, E_2,$ and $E_3$ when $t = 0.1, x = 0.2, \varepsilon = 0.2, \phi = 0.1, M = 1.0, BrA^{-1} = 1.0, Gr = 0.03, Br = 3.0, \beta_1 = 0.01, \beta_2 = -0.01, \gamma_1 = 0.01, \text{and } \gamma_2 = -0.01.$

Figures 19–23 demonstrate the behavior of Bejan number with pertinent parameters. Figure 19 displays the impact of nanoparticles volume fraction on Be. The relation between Bejan number and nanoparticles volume fraction is inverse as increase in $\phi(0.1 – 0.7)$
Figure 19. Be in terms of $\phi$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $M = 1.0$, $\text{BrA}^{-1} = 1.0$, $\text{Gr} = 0.03$, $\text{Br} = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 20. Be in terms of $M$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $\phi = 0.1$, $\text{BrA}^{-1} = 1.0$, $\text{Gr} = 0.03$, $\text{Br} = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 21. Be in terms of $\text{Gr}$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $\phi = 0.1$, $M = 1.0$, $\text{BrA}^{-1} = 1.0$, $\text{Br} = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 22. Be in terms of $\text{BrA}^{-1}$ when $t = 0.1$, $E_3 = 0.01$, $x = 0.2$, $E_2 = 0.01$, $E_1 = 0.02$, $\varepsilon = 0.2$, $\phi = 0.1$, $M = 1.0$, $\text{Gr} = 0.03$, $\text{Br} = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Figure 23. Be in terms of $E_1$, $E_2$, and $E_3$ when $t = 0.1$, $x = 0.2$, $\varepsilon = 0.2$, $\phi = 0.1$, $M = 1.0$, $\text{BrA}^{-1} = 1.0$, $\text{Gr} = 0.03$, $\text{Br} = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$.

Bejan number increases with rise in ratio of Br to A from 0.1 to 0.7. Figure 23 presents the results for wall parameters. While enhancement is observed in Bejan number with larger elastance parameters, i.e. $E_1 (= 0.01, 0.02)$ and $E_2 (= 0.02, 0.04)$, it is decreased with larger values of the damping parameter, i.e. $E_3 (= 0.01, 0.02)$. Moreover, in all cases, the values of the parameters in Hamilton-Cropper’s model are less than those in Maxwell’s model.

4.4. Streamlines
The streamlines are plotted to illustrate trapping. Figure 24(a) and (b) shows the impact of Hartman number in Maxwell model and Figure 24(c) and (d) portrays it in Hamilton-Cropper model. In both cases, the size of trapped bulus increases with larger values of Hartman number ($M = 1.0, 2.0$). Figures 25 and 26(a)-(d) picture the behavior of the first- and second-order slip parameters. The streamlines indicate that the size of the trapped bulus rises with enhancement in the first-order slip (from 0.01 to 0.03) and second-order slip parameter (−0.01, −0.03). The impact of wall parameters can be observed in Figure 27(a)-(d).
Figure 24. $\psi$ in terms of $M$ for (a) and (b) Maxwell model and (c) and (d) Hamilton-Crosser model when $E_3 = 0.01$, $t = 0$, $E_2 = 0.01$, $\varepsilon = 0.2$, $E_1 = 0.02$, $\phi = 0.1$, $Gr = 0.03$, $Br = 3.0$, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$: (a) and (c) $M = 1.0$ and (b) and (d) $M = 2.0$.

Figure 25. $\psi$ in terms of $\beta_1$ for (a) and (b) Maxwell model and (c) and (d) Hamilton-Crosser model when $E_3 = 0.01$, $t = 0$, $E_2 = 0.01$, $\varepsilon = 0.2$, $E_1 = 0.02$, $\phi = 0.1$, $M = 1.0$, $Gr = 0.03$, $Br = 3.0$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$: (a) and (c) $\beta_1 = 0.01$, and (b) and (d) $\beta_1 = 0.03$

for the Maxwell model and in Figure 27(e)-(h) for the Hamilton-Crosser model. The models show the same behavior with these parameters, i.e. trapped bolus size increases with $E_1(=0.7,0.9)$ and $E_2(=0.4,0.6)$ and decreases with $E_3(=0.2,0.5)$.

5. Conclusions

The main conclusions drawn from the observations in this study may be summarized as follows:

- Enhancement in velocity was observed with both first-order and second-order velocity slips in Maxwell and Hamilton-Crosser models, while it was reduced with nanoparticle volume fraction;
- Values of the parameters were higher in Hamilton-Crosser model than those in Maxwell model, especially near the center of the channel for velocity profile;
- Grashof and Hartman numbers had reverse effect on velocity;
- The quantities of temperature in Maxwell model exceeded those in Hamilton-Crosser model;
- Larger values of $Br\Lambda^{-1}$ and Grashof number led to increase in entropy generation. On the other hand, Hartman number and nanoparticle volume fraction had an inverse impact on entropy generation;
- Bolus sizes increased by trapping phenomenon with both first- and second-order velocity slip parameters;
Figure 27. $\psi$ in terms of $E_1$, $E_2$, and $E_3$ for (a), (b), (c) and (d) Maxwell model and (e), (f), (g) and (h) Hamilton-Crosser model when $t = 0$, $\varepsilon = 0.2$, $\phi = 0.1$, $M = 1.0$, Gr = 0.03, Br = 3.0, $\beta_1 = 0.01$, $\beta_2 = -0.01$, $\gamma_1 = 0.01$, and $\gamma_2 = -0.01$: (a) and (e) $E_1 = 0.7$, $E_2 = 0.4$, and $E_3 = 0.2$; (b) and (f) $E_1 = 0.9$, $E_2 = 0.4$, and $E_3 = 0.2$; (c) and (g) $E_1 = 0.7$, $E_2 = 0.6$, $E_3 = 0.2$; (d) and (h) $E_1 = 0.7$, $E_2 = 0.4$, and $E_3 = 0.5$.

- Bolus sizes were reduced with $E_3$ and increased with $E_1$ and $E_2$ in both Maxwell and Hamilton-Crosser models.

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