Measurements of branching fraction ratios and CP asymmetries in $B^{\pm} \to D_{CP}K^{\pm}$ decays in hadron collisions

T. Aaltonen,24 J. Adelman,14 B. Álvarez González,12 S. Amerio,44 D. Amidei,35 A. Anastassov,39 A. Anovn,20 J. Antos,15 G. Apollinari,18 A. Apresyan,49 T. Arisawa,38 A. Artikov,16 J. Asadi,54 W. Ashmanskas,18 A. Attal,4 A. Aurisano,54 F. Azfar,33 W. Badgett,18 A. Barbaro-Galtieri,29 V.E. Barnes 49 B.A. Barnett,26 P. Barria,47 P. Bartos,15 G. Bauer,33 P.-H. Beauchemin,34 F. Bedeschi,47 D. Beecher,31 S. Behari,26 G. Bellettini,47 J. Bellinger,60 D. Benjamine,17 A. Berervas,18 A. Bhatti,51 M. Binkley,18 D. Bisello,44 I. Bizzaj,31 R.E. Blair,2 C. Blocker,7 B. Blumenfeld,26 A. Bocci 17 A. Bodek,50 V. Boisvert,50 D. Bortoletto,49 J. Bondreau,48 A. Boveia,11 B. Brau,14 A. Bridgeman,25 L. Brigliadori,6 C. Bromberg,36 E. Brubaker,14 J. Budagov,16 H.S. Budd,50 S. Budd,25 K. Burkett,18 G. Busetto,44 P. Bussey,22 A. Buzatu,24 K. Byrum,2 S. Cabrera,17 C. Calanca,32 S. Camarda,4 M. Campanelli,36 M. Campbell,35 F. Canela,14 A. Canepa,46 B. Carls,25 D. Carlsmith,60 R. Carosi,47 S. Carrillo,19 S. Carron,18 B. Casal,12 M. Casarsa,18 A. Castro,6 C. Catastini,47 D. Cauz,55 V. Cavaliere,47 M. Cavalli-Sforza,4 A. Cerri,29 L. Cerrito,31 S.H. Chang,28 Y.C. Chen,1 M. Chertok,8 G. Chiarelli,47 G. Chlachidze 18 F. Chlebana,18 K. Cho,28 D. Chokheli,16 J.P. Chou,23 K. Chung,18 W.H. Chung,60 Y.S. Chung,50 T. Chwalek,27 C.I. Ciobanu,45 M.A. Ciocci,47 A. Clark,21 D. Clark,7 G. Compostella,44 M.E. Convery,17 J. Conway,6 M. Corbo,45 M. Cordelli,26 C.A. Cox,6 D.J. Cox,6 F. Crescioli,47 C. Cuenca Almenar,61 J. Cuevas,12 R. Culbertson,18 J.C. Cully,35 D. Dagenhart,18 M. Datta,18 T. Davies,22 P. de Barbaro,50 S. De Cecco,52 A. Deisher,29 G. De Lorenzo,4 M. Dell’Orso,47 C. Deluca,4 L. Demortier,51 J. Deng,17 M. Denino,6 M. Derriero,44 A. Di Canto,47 G.P. di Giovanni,45 B. Di Ruzza,47 J.R. Dittmann,5 M. D’Onofrio,47 S. Donati,47 P. Dong,18 T. Dorigo,44 S. Dube,53 K. Ebina,58 A. Elagin,54 R. Erbacher,8 D. Errede,25 S. Errede,25 N. Esralettah,45 R. Eusebi,54 H.C. Fang,29 S. Farrington,43 W.T. Fedorko,14 R.G. Feild,61 M. Feindt,27 J.P. Fernandez,32 C. Ferrazza,49 R. Field,9 G. Flanagan,9 R. Forrest,8 M.J. Frank,5,5 Franklin,23 J.C. Freeman,18 I. Furic,19 M. Gallinaro,51 J. Galyardt,13 F. Garbers,11 J.E. Garcia,21 A.F. Garfinkel,9 P. Garosi,47 H. Gerberich,29 D. Gerdes,35 A. Gessler,27 S. Giagu,52 V. Giakoumopoulou,3 P. Giannetti,47 K. Gibson,47 J.L. Gimmel,50 C.M. Ginsburg,18 N. Giokarisi,3 M. Giordani,54 P. Girgiani,20 M. Giunta,47 G. Giurgiu,26 V. Glagolev,16 M. Gold,38 N. Golshmidt,19 A. Golossanov,18 G. Gomez,12 Gomez-Caballlos,33 M. Goncharov,33 O. Gonzalez,32 I. Gorenko,38 A.T. Goshaw,17 K. Goulianos,51 A. Greisele,44 S. Grinstein,5 C. Grosso-Pilcher,14 R.C. Group,18 U. Grundler,25 J. Guimaraes da Costa,29 Z. Gunay-Unalan,36 C. Haber,29 S.R. Han,18 E. Halkiadakis,53 B.-Y. Han,50 J.Y. Han,50 F. Happacher,12 K. Hara,56 D. Hare,53 M. Hare,57 R.F. Harr,59 M. Hartz,48 K. Hatakeyama,5 C. Hays,43 M. Heck,27 J. Heinrich,46 M. Herndon,50 J. Heuser,27 S. Hewamanage,5 D. Hidas,53 C.S. Hill,11 D. Hirschbuehl,27 A. Hocker,18 S. Hou,1 M. Houliden,30 S.-C. Hsu,29 R.E. Hughes,40 M. Hurwitz,14 U. Husemann,61 M. Hussein,36 J. Huston,36 J. Incandela,11 G. Introzzi,47 M. Iori,52 A. Ivanov,8 E. James,18 D. Jang,13 B. Jayatilaka,17 E.J. Jeon,28 M.K. Jha,6 S. Jindariani,18 W. Johnson,5 M. Jones,49 K.K. Joo,26 S.Y. Jun,13 J.E. Jung,28 T.R. Junk,18 T. Kamon,24 D. Kaur,19 P.E. Karchin,59 Y. Kato,90,42 R. Kepl,18 W. Ketchum,14 J. Keung,46 V. Khotilovich,54 B. Kilminster,18 D.H. Kim,28 H.S. Kim,28 H.W. Kim,28 J.E. Kim,26 M.J. Kim,28 S.B. Kim,28 S.H. Kim,56 Y.K. Kim,14 N. Kimura,58 L. Kirsch,7 S. Klimenko,19 K. Kondo,56 D.J. Kong,28 J. Konigsberg,19 A. Korytov,19 A. Kowalski,17 M. Krep,27 J. Kroll,46 D. Krop,14 N. Krumnack,5 M. Kruse,17 V. Krutelov,11 T. Kuhl,27 N.P. Kulkarni,59 M. Kurata,56 S. Kwang,14 A.T. Laasanen,49 S. Lami,47 S. Lammel,18 M. Lancaster,31 R.L. Landier,8 K. Lannon,40 A. Lath,53 G. Latino,47 I. Lazzizzera,44 T. LeCompte,2 E. Lee,54 H.S. Lee,14 J.S. Lee,28 S.W. Lee,54 S. Leone,47 J.D. Lewis,18 C.-J. Lin,29 J. Linacre,43 M. Lindgren,18 E. Lipeles,46 A. Lister,21 D.O. Litvinsev,18 C. Liu,48 T. Liu,18 N.S. Lockyer,14 A. Logino,61 L. Lovas,15 D. Lucchesi,44 J. Lueck,27 P. Lujan,29 P. Lukens,18 G. Lungu,51 J. Lys,29 R. Lysak,34 R. Madrak,18 K. Maeshima,18 K. Makhoul,33 P. Maksimovic,26 S. Malde,43 S. Malik,31 G. Manca,30 A. Manousakis-Katsikakis,3 F. Margaroli,49 C. Marino,27 C.P. Marino,25 A. Martin,61 V. Martín,24 M. Martínez,4 R. Martínez-Ballarín,32 P. Mastrandrea,52 M. Mathis,26 M.E. Mattson,5 P. Mazzanti,6 K.S. McFarland,50 P. McIntyre,54 R. McNulty,30 A. Melha,50 P. Melhata,24 A. Menzione,47 C. Mesropian,51 T. Miao,18 D. Mietlicki,35 N. Miladinovic,7 R. Miller,36 C. Mills,23 M. Milnik,27 A. Mitra,51 G. Mitselmakher,19 H. Miyake,36 S. Moed,23 N. Moggi,6 M.N. Mondragon,28 C.S. Moon,28 R. Moore,18 M.J. Morello,47 J. Morlock,27 M.P. Movila Fernandez,18 J. Müllennstädt,29 A. Mukherjee,18 Th. Mulder,27 P. Murat,18 M. Mussinevic,6 J. Nachtman,18 Y. Nagai,56 J. Nagano,56 K. Nakamura,56 I. Nakano,41 A. Napier,57 J. Nett,60 C. Neu,46 M.S. Neubauer,25 S. Neubauer,27 J. Nielsen,929 L. Nodulman,2 M. Norman,10 O. Noernaell,25 E. Nurse,31 L. Oakes,43 S.H. Oh,37 Y.D. Oh,28 I. Oksuzian,19 T. Okusawa,42 R. Orava,24 K. Osterberg,24 S. Pagan Griso,44
We reconstruct $B^\pm \rightarrow DK^\pm$ decays in a data sample collected by the CDF II detector at the Tevatron collider corresponding to 1 fb$^{-1}$ of integrated luminosity. We select decay modes where the $D$ meson decays to either $K^-\pi^+$ (flavor eigenstate) or $K^-\pi^+\pi^-\pi^+$ ($CP$-even eigenstates), and measure the direct $CP$ asymmetry $A_{CP^+} = 0.39 \pm 0.17$ (stat) $\pm 0.04$ (syst), and the double ratio of $CP$-even to flavor eigenstate branching fractions $R_{CP^+} = 1.30 \pm 0.24$ (stat) $\pm 0.12$ (syst). These measurements will improve the determination of the Cabibbo-Kobayashi-Maskawa angle $\gamma$. They are performed here for the first time using data from hadron collisions.

PACS numbers: 13.25.Hw 11.30.Er 14.40.Nd
The measurement of CP asymmetries and branching ratios of $B^- \rightarrow DK^-$ decay modes allows a theoretically-clean extraction of the CKM angle $\gamma = \arg(-V_{ub} V_{us}^*/V_{cd} V_{cs}^*)$, a fundamental parameter of the standard model [2]. In these decays the interference between the tree amplitudes of the $b \rightarrow c\bar{s}s$ and $b \rightarrow u\bar{c}c$ processes leads to observables that depend on their relative weak phase ($\gamma$), their relative strong phase ($\delta_B$), and the magnitude ratio $R_B = \frac{|A_{b\rightarrow c\bar{s}s}|}{|A_{b\rightarrow u\bar{c}c}|}$. These quantities can all be extracted from data by combining several experimental observables. This can be achieved in several ways, from a variety of $D$ decay channels [3, 5].

An accurate knowledge of the value of $\gamma$ is instrumental in establishing the possible presence of additional non-standard model CP-violating phases in higher-order diagrams [3, 5]. Its current determination is based on a combination of several $B \rightarrow DK$ measurements performed in $e^+e^-$ collisions at the T(4S) resonance [8] and its uncertainty is between 12 and 30 deg, depending on the method [11]. This uncertainty is almost completely determined by the limited size of the data samples available, with theoretical uncertainties playing a negligible role ($\sim 1\%$). The large production of $B$ mesons available at hadron colliders could offer a unique opportunity to improve the current experimental determination of the angle $\gamma$. However, the feasibility of this kind of measurement in the larger background conditions of hadronic collisions has never been demonstrated.

In this paper we describe the first measurement of the branching fraction ratios and CP asymmetries of $B^- \rightarrow DK^-$ modes performed in hadron collisions, based on an integrated luminosity of 1 fb$^{-1}$ of $\bar{p}p$ collisions at $\sqrt{s} = 1.96$ TeV collected by the upgraded Collider Detector (CDF II) at the Fermilab Tevatron. We reconstruct events where the $D$ meson decays to the flavor-specific mode $K^-\pi^+$ ($D^0$), or to one of the CP-even modes $K^-K^+$ and $\pi^-\pi^+$ [$D_{CP^+} = (D^0 + \overline{D}^0)/\sqrt{2}$]. From these modes, the following observables can be defined:

$$ A_{CP^+} = \frac{B(B^- \rightarrow D_{CP^+}K^-) - B(B^+ \rightarrow D_{CP^+}K^+)}{B(B^- \rightarrow D_{CP^+}K^-) + B(B^+ \rightarrow D_{CP^+}K^+)}. $$

$$ R_{CP^+} = \frac{B(B^- \rightarrow D_{CP^+}K^-) + B(B^+ \rightarrow D_{CP^+}K^+)}{B(B^- \rightarrow D^0_{CP^+}K^-) + B(B^+ \rightarrow \overline{D}^0_{CP^+}K^+)} $$

With the assumption of no CP violation in $D^0$ decays, and neglecting $D^0-\overline{D}^0$ mixing [12], these quantities are related to the CKM angle $\gamma$ by the equations [3]

$$ R_{CP^+} = 1 + 2\gamma^2 + 2r\cos\delta_B\cos\gamma, $$

$$ A_{CP^+} = 2r_B\sin\delta_B\sin\gamma/R_{CP^+}. $$

For our measurements we adopt the usual approximation $R_{CP^+} \sim \frac{R_B}{R_{13}}$, which is valid up to a term $r \cdot |V_{us} V_{cd}/V_{ud} V_{cs}| \approx 0.01_{13}$, where

$$ R = \frac{B(B^- \rightarrow D^0_{CP^+}K^-) + B(B^+ \rightarrow \overline{D}^0_{CP^+}K^+)}{B(B^- \rightarrow D^0_{CP^+}\pi^-) + B(B^+ \rightarrow \overline{D}^0_{CP^+}\pi^+)}, $$

$$ R_+ = \frac{B(B^- \rightarrow D_{CP^+}K^-) + B(B^+ \rightarrow D_{CP^+}K^+)}{B(B^- \rightarrow D_{CP^+}\pi^-) + B(B^+ \rightarrow D_{CP^+}\pi^+)} $$

The CDF II detector is a multipurpose magnetic spectrometer surrounded by calorimeters and muon detectors. The components relevant for this analysis are briefly described here. A more detailed description can be found elsewhere [13]. Silicon microstrip detectors (SVX II and ISL) [14] and a cylindrical drift chamber (COT) [15] immersed in a 1.4 T solenoidal magnetic field allow reconstruction of charged particles in the pseudorapidity range $|\eta| < 1.0$ [16]. The SVX II detector consists of microstrip sensors arranged in five concentric layers with radii between 2.5 and 10.6 cm, divided into three contiguous sections along the beam direction $z$, for a total length of 90 cm. The two additional silicon layers of the ISL help to link tracks in the COT to hits in the SVX II. The COT has 96 measurement layers between 40 and 137 cm in radius, organized into alternating axial and $\pm 2\theta$ stereo superlayers, and provides a resolution on the transverse momentum of charged particles $p_T \approx 0.15p_T/(GeV/c)$. The specific energy loss by ionization ($dE/dx$) of charged particles in the COT can be measured from the collected charge, which is encoded in the output pulse width of each sense wire.

Candidate events for this analysis are selected by a three-level trigger system. At level 1, charged particles are reconstructed in the COT axial superlayers by a hardware processor, the extremely fast tracker (XFT) [16]. Two oppositely charged particles are required, with transverse momenta $p_T \geq 2$ GeV/c and scalar sum $p_{T1} + p_{T2} \geq 5.5$ GeV/c. At level 2, the silicon vertex trigger (SVT) [17] associates SVX II $r-\phi$ position measurements with XFT tracks. This provides a precise measurement of the track impact parameter, $d_0$, which is defined as the distance of closest approach to the beam line. The resolution of the impact parameter measurement is 50 pm for particles with $p_T$ of about 2 GeV/c, including a $\sim 30$ pm contribution due to the transverse beam size, and improves for higher transverse momenta. We select $B$ hadron candidates by requiring two SVT
tracks with 120 $\leq d_{0} \leq 1000 \mu m$. To reduce background from light-quark jet pairs, the two trigger tracks are required to have an opening angle in the transverse plane $2^\circ \leq \Delta\phi \leq 90^\circ$, and to satisfy the requirement $L_{xy} > 200 \mu m$, where $L_{xy}$ is defined as the distance in the transverse plane from the beam line to the two-track vertex, projected onto the two-track momentum vector. The level 1 tracks with $120 < p_{T} < 0.4$ GeV to form $B^{-}$ candidates. A kinematic fit of the decay is performed by constraining the two tracks forming the $D^{0}$ candidate to a common vertex and to the nominal $D^{0}$ mass, the $D^{0}$ candidate and the remaining track to a separate vertex, and the reconstructed momentum of the $B^{-}$ candidate to point back to the luminous region in the transverse plane.

To complete the selection, further requirements are applied on additional observables: the impact parameter ($d_{0}$) of the reconstructed $B$ candidate relative to the beamline; the isolation of the $B$ candidate ($I_{B}$); the goodness of fit of the decay vertex ($\chi_{B}^{2}$); the transverse distance of the $D$, both relative to the beam $[L_{xy}(D)]$ and to the $B$ vertex $[L_{xyB}(D)]$, and the significance of the $B$ hadron decay length $[L_{xy}(B)/\sigma _{L_{xy}(B)}]$. We chose the requirement $L_{xyB}(D) > 100 \mu m$ to reduce contamination from (nonresonant) three-body decays of the type $B^{+} \rightarrow h^{+}h^{-}h^{+}$ (from here on, we will use $h$ to indicate either $K$ or $\pi$), in which all tracks come from a common decay vertex. In addition, we reject all candidates comprising a pair of tracks with an invariant mass compatible with a $J/\psi \rightarrow \mu^{+}\mu^{-}$ decay within $2\sigma$. The threshold values for all other requirements, whose purpose is to reduce combinatorial background, were determined by an unbiased optimization procedure aimed at achieving the best resolution on $A_{CP}$. This resolution was parametrized as a function of the expected signal yield $S$ and background level $B$, by performing repeated fits on samples of simulated data extracted from the same multidimensional distribution used as likelihood function in the fit [Eq. (7)]. For each choice of thresholds, the signal $S$ was determined by rescaling the number of observed $B^{-} \rightarrow D_{s}^{0} \pi^{-}$, and the background $B$ was determined from the upper mass sidebands of each data sample ($5.4 < M_{B} < 5.8$ GeV/$c^{2}$). Based on this optimization procedure, we adopted the following set of requirements: $I_{B} > 0.65$, $\chi_{B}^{2} < 13$, $d_{0} < 70 \mu m$, $L_{xy}(B)/\sigma _{L_{xy}(B)} > 12$, and $L_{xy}(D) > 400 \mu m$.

For every $B^{-} \rightarrow Dh^{-}$ candidate, a nominal invariant mass is evaluated by assigning the charged pion mass to the particle $h^{-}$ coming from the $B$ decay. The distributions obtained for the three modes of interest ($D \rightarrow K\pi$,$KK$ or $\pi\pi$) are reported in Fig. 1. A clear $B^{-} \rightarrow D_{s}\pi^{-}$ signal is seen in each. Events from $B^{-} \rightarrow DK^{-}$ decays are expected to form much smaller and wider peaks in these plots, located about 50 MeV/$c^{2}$ below the $B^{-} \rightarrow D\pi^{-}$ peaks, and as such cannot be resolved. The dominant residual backgrounds are random track combinations that meet the selection requirements (combinatorial background), misreconstructed physics background such as $B^{-} \rightarrow D^{*0}\pi^{-}$ decay, and, in the $D^{0} \rightarrow KK$ final state, the nonresonant $B^{-} \rightarrow K^{+}K^{-}K^{-}$ decay, as determined by a study performed on CDF simulation.

We used an unbinned likelihood fit, exploiting kinematic and particle identification information from the measurement of $dE/dx$ in a similar way to $[21]$, to separate statistically the $B^{-} \rightarrow DK^{-}$ contributions from the $B^{-} \rightarrow D\pi^{-}$ signals and from the combinatorial background. To make best use of the available information, we fit the three modes simultaneously using a single likelihood function, to take advantage of the presence of parameters common to the three modes.

The likelihood function is

$$ L = \prod_{i=1}^{n} \left[ (1-b) \sum_{j} f_{j} \mathcal{L}_{j}^{\text{kin}} \mathcal{L}_{j}^{\text{PID}} + b \mathcal{L}_{c}^{\text{kin}} \mathcal{L}_{c}^{\text{PID}} \right] \tag{7} $$

where $c$ labels combinatorial background quantities, $b$ is the combinatorial background fraction, and $\mathcal{L}_{j}^{\text{kin}}$ and $\mathcal{L}_{j}^{\text{PID}}$ are defined below. The index $j$ runs over the modes $B^{-} \rightarrow DK^{-}$, $B^{-} \rightarrow D\pi^{-}$, nonresonant $B^{-} \rightarrow K^{+}K^{-}K^{-}$ and $B^{-} \rightarrow \pi^{+}\pi^{-}K^{-}$, and $B^{-} \rightarrow D^{*0}\pi^{-}$ (where a soft $\gamma$ or $\pi^{0}$ from the $D^{*0}$ is undetected) and $f_{j}$ are the fractions to be determined by the fit. The fraction of the physics background ($B^{-} \rightarrow D^{*0}\pi^{-}$) with respect to the signal is common to the three decays and the fraction of the $B^{-} \rightarrow D_{CP}^{+}\pi^{-}$ is common to the two $D_{CP}$ modes. As determined from simulation, these modes are the only significant contributions within the mass range $5.17 < M < 5.60$ GeV/$c^{2}$ chosen for our fit.

Kinematic information is given by three loosely correlated observables: (a) the mass $M_{D_{s}}$, calculated by assigning the pion mass to the track from the $B$ decay; (b) the momentum imbalance $\alpha$, defined as

$$ \alpha = 1 - p_{tr}/p_{D} > 0 \quad \text{if} \quad p_{tr} < p_{D}; \quad \alpha = -(1 - p_{D}/p_{tr}) \leq 0 \quad \text{if} \quad p_{tr} \geq p_{D}; $$

where $p_{tr}$ is the momentum of the track from the $B$ candidate; and (c) the scalar sum of the $D$ momentum and the momentum of the track from the $B$ candidate ($p_{tot} = p_{tr} + p_{D}$). The above variables uniquely identify the invariant mass $M_{DK}$ evaluated with a kaon mass assignment to the track from the $B$ decay, through the (exact) relations

$$\chi_{B}^{2} < 13, \quad d_{0} < 70 \mu m, \quad L_{xy}(B)/\sigma _{L_{xy}(B)} > 12, \quad L_{xy}(D) > 400 \mu m.$$
FIG. 1: Invariant mass distributions of $B^- \rightarrow D\pi^-$ candidates for each reconstructed decay mode. The pion mass is assigned to the charged track from the $B$ candidate decay vertex. The projections of the common likelihood fit (see text) are overlaid for each mode.

$$M^2_{DK} = M_{D\pi}^2 + m_{\pi}^2 - m_K^2 + 2\sqrt{m_D^2 + \frac{p_{tot}^2}{(2-\alpha)^2}} \left( \sqrt{m_D^2 + \frac{(p_{tot}(1-\alpha))^2}{2-\alpha}} - \sqrt{m_K^2 + \frac{(p_{tot}(1-\alpha))^2}{2-\alpha}} \right)$$

if $\alpha > 0$;

$$M^2_{DK} = M_{D\pi}^2 + m_{\pi}^2 - m_K^2 + 2\sqrt{m_D^2 + \frac{p_{tot}(1+\alpha)^2}{2+\alpha}} \left( \sqrt{m_D^2 + \frac{(p_{tot}(1+\alpha))^2}{2+\alpha}} - \sqrt{m_K^2 + \frac{(p_{tot}(1+\alpha))^2}{2+\alpha}} \right)$$

if $\alpha \leq 0$.

Using these variables, we can write $L^{{\text{kin}}}_j = P_j(M_{D\pi}|\alpha, p_{tot})P_j(\alpha, p_{tot})$ and $L^{{\text{PID}}}_j = P_j(dE/dx|\alpha, p_{tot})$, where $P_j$ is the probability density function for decay mode $j$. Distributions of the kinematic variables for the signals are obtained from samples of events from the full CDF simulation, while the combinatorial background they are obtained from the mass sidebands of data. The shape of the mass distribution assigned to each signal process ($B^- \rightarrow D\pi^-$ and $B^- \rightarrow DK^-$ decays) has been modeled in detail from a dedicated study including the effect of final state QED radiation [22]. The simulation results were tested on high-statistics data samples of $D^0$ decays, in order to ensure the reliability of the extraction of the $DK^-$ component in the vicinity of the larger $D\pi^-$ peak. Exponential functions were used to model the mass distribution of combinatorial background for each mode. The normalizations and the slope of these functions are independently determined in the maximum likelihood fit. The particle identification (PID) model of the combinatorial background allows for pion and kaon components, which are free to vary in the fit.

A large sample of $D^{*+} \rightarrow D^0(\rightarrow K^-\pi^+)\pi^+$ decays was used to calibrate the $dE/dx$ response of the detector to kaons and pions, using the charge of the pion in the $D^{*+}$ decay to determine the identity of the $D^0$ decay products. The calibration includes the dependence of the shape and the average of the response curve on particle momentum, and the shape of the distribution of common-mode fluctuations. The calibrated $dE/dx$ information provides a 1.5 $\sigma$ separation power between pion and kaon particles of $p_T > 2$ GeV/c. Uncertainties on the calibration parameters are included in the final systematic uncertainty of $A_{CP^+}$ and $R_{CP^+}$ [22].

The $B^- \rightarrow DK^-$ and $B^- \rightarrow D\pi^-$ signal event yields obtained from the fit to the data are reported in Table. The fraction of the $B^- \rightarrow \pi^+\pi^-K^-$ was set by the fit to its lower bound at zero, compatible with the expectation of a negligible contribution, and will be ignored in the following. The uncorrected values of the double ratio of branching fractions $R_{CP^+}$ and of the CP asymmetry $A_{CP^+}$ obtained from the fit are $R_{CP^+} = 1.27 \pm 0.24$ and $A_{CP^+} = 0.39 \pm 0.17$. In the fit, $R_{CP^+}$ and $A_{CP^+}$ are functions of the fractions $f_j$ in Eq. [4] and the total number of events in each subsample.

As a check of the goodness of the fit, and to visualize better the separation between signal and background, we
plot distributions of the relative signal likelihoods:

\[
R_{\text{L}} = \frac{pdf(B \to DK)}{pdf(B \to DK) + pdf(\text{background})}
\]

where \(pdf(B \to DK)\) is the probability density under the signal hypothesis, and \(pdf(\text{background})\) is the probability density under the background hypothesis (including both physics and combinatorial backgrounds, with their measured relative fractions). These distributions are compared to the prediction of our fit in Fig. 2 showing a very good agreement. In addition, we plot projections of the fit on the invariant mass distributions, both for the entire sample (Fig. 1), and for a kaon–enriched subsample, where the interesting \(B^- \to DK^-\) components have been enhanced with respect to the \(B^- \to D\pi^-\) by means of a \(dE/dx\) cut (Fig. 3). All these projections show very good agreement between our fit and the data.

Some corrections are needed to convert our fit results into measurements of the parameters of interest. First, we correct for small biases in the fit procedure itself, as measured by repeated fits on simulated samples: \(\delta(R_{\text{CP}^{+}}) = -0.027 \pm 0.005\) and \(\delta(A_{\text{CP}^{+}}) = 0.015 \pm 0.003\). These biases are independent of the true values of \(A_{\text{CP}^{+}}\) and \(R_{\text{CP}^{+}}\) used in the simulated samples. \(R_{\text{CP}^{+}}\) does not need any further corrections because detector effects cancel in the double ratio of branching fractions. The direct \(CP\) asymmetry \(A_{\text{CP}^{+}}\) needs to be corrected for the different probability for \(K^+\) and \(K^-\) mesons to interact with the tracker material. This effect is reproduced well by CDF II detector simulation (traced by GEANT [24]), which yields an estimate \(\frac{\sigma(K^+)}{\sigma(K^-)} = 1.0178 \pm 0.0023\) (stat) \(\pm 0.0045\) (syst) [25] which has been verified by measurements on data [20].

The corrected results are

\[
R_{\text{CP}^{+}} = 1.30 \pm 0.24\text{(stat)}, \quad A_{\text{CP}^{+}} = 0.39 \pm 0.17\text{(stat)},
\]

where \(A_{\text{CP}^{+}}\) was corrected using the following equation:

\[
A_{\text{CP}^{+}} = \frac{N(B^- \to D_{\text{CP}^{+}}^0K^-) \frac{\sigma(K^+)}{\sigma(K^-)} - N(B^+ \to D_{\text{CP}^{+}}^0K^+)}{N(B^- \to D_{\text{CP}^{+}}^0K^-) \frac{\sigma(K^+)}{\sigma(K^-)} + N(B^+ \to D_{\text{CP}^{+}}^0K^+)},
\]

Systematic uncertainties are listed in Table II. They were determined by generating simulated samples of pseudoexperiments with different underlying assumptions, and checking the effect of such changes on the results of our measurement procedure. The dominant contributions are uncertainty on the \(dE/dx\) calibration and parametrization, uncertainty on the kinematics of the combinatorial background, and uncertainty on the physics background \((B^- \to D^{\ast\ast}\pi^-)\) mass distribution. Variations in the model of the combinatorial background included different functional forms of the mass distribution, and alternative \((\alpha, p_{\text{cut}})\) distributions, constrained by comparison with real data in the mass sidebands.

Smaller contributions are assigned for trigger efficiencies, assumed \(B^-\) mass input in the fit [27] and kinematic properties of signal and physics background.

In summary, we have measured the double ratio of \(CP\) to flavor eigenstate branching fractions [Eq. 2] \(R_{\text{CP}^{+}} = 1.30 \pm 0.24\) (stat) \(\pm 0.12\) (syst) and the direct \(CP\) asymmetry [Eq. 1] \(A_{\text{CP}^{+}} = 0.39 \pm 0.17\) (stat) \(\pm 0.04\) (syst). These results can be combined with other \(B^- \to DK^-\) decay parameters to improve the determination of the CKM angle \(\gamma\). These measurements are performed here for the first time in hadron collisions, are in agreement with previous measurements from BaBar \((R_{\text{CP}^{+}} = 1.06 \pm 0.10 \pm 0.05, A_{\text{CP}^{+}} = 0.27 \pm 0.09 \pm 0.04\) in 348 fb\(^{-1}\) of integrated luminosity [3]) and Belle \((R_{\text{CP}^{+}} = 1.13 \pm 0.16 \pm 0.08, A_{\text{CP}^{+}} = 0.06 \pm 0.14 \pm 0.05\) in 250 fb\(^{-1}\) of integrated luminosity [10]) and have comparable uncertainties.

\[
\begin{array}{cccc}
\hline
D \text{ mode} & B^+ \to D\pi^+ & B^- \to D\pi^- & B^- \to DK^- & B^- \to D\pi^- \to h^- h^+ \to h^- h^+ \to h^- h^+ \\
\hline
K^\pm \pi^\mp & 3769 \pm 68 & 3763 \pm 68 & 250 \pm 26 & 266 \pm 27 & - & - \\
K^+ K^- & 381 \pm 25 & 399 \pm 26 & 22 \pm 8 & 49 \pm 11 & 3 \pm 1 & 3 \pm 1 \\
\pi^+ \pi^- & 101 \pm 13 & 117 \pm 14 & 6 \pm 6 & 14 \pm 6 & - & - \\
\hline
\end{array}
\]

| Source | \(R_{\text{CP}^{+}}\) | \(A_{\text{CP}^{+}}\) |
|--------|-----------------|----------------|
| \(dE/dx\) model | 0.056 | 0.030 |
| \(D^{\ast\ast}\pi\) mass model | 0.025 | 0.006 |
| Input \(B^-\) mass to the fit | 0.004 | 0.002 |
| Combinatorial background mass model | 0.020 | 0.001 |
| Combinatorial background kinematics | 0.100 | 0.020 |
| \(D\pi\) kinematics | 0.002 | 0.001 |
| \(DK\) kinematics | 0.002 | 0.004 |
| \(D^{\ast\ast}\pi\) kinematics | 0.004 | 0.002 |
| Fit bias | 0.005 | 0.003 |
| Total (sum in quadrature) | 0.12 | 0.04 |

TABLE I: \(B^- \to DK^-\) and \(B^- \to D\pi^-\) event yields obtained from the fit to the data.
FIG. 2: Relative likelihood for $B^- \to DK^-$ candidates for each reconstructed decay mode. The points with the error bars show the distribution obtained on the fitted data sample while the histograms show the distributions obtained by generating signal and background events directly from the total PDF of the fit composition.

FIG. 3: Invariant mass distributions of $B^- \to D\pi^-$ candidates for each reconstructed decay mode. The pion mass is assigned to the prompt track from the $B$ decay. A requirement on the PID variable was applied to suppress the $D\pi$ component and favor the $DK$ component. The projections of the likelihood fit for each mode are overlaid. The $p$-value for agreement of data with the fit is 0.95

Acknowledgments

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research
Council of Canada; the National Science Council of the Republic of China; the A.P. Sloan Foundation; the Bundesministerium für Bildung und Forschung, Germany; the Korean Science and Engineering Foundation and the Korean Research Foundation; the Science and Technology Facilities Council and the Royal Society, UK; the Institut National de Physique Nucléaire et Physique des Particules/CNRS; the Russian Foundation for Basic Research; the Ministerio de Ciencia e Innovación, and Programa Consolider-Ingenio 2010, Spain; the Slovak R&D Agency; and the Academy of Finland.

[1] The charge conjugate state is implied throughout the paper, except in formulas and sentences where both are mentioned explicitly.

[2] M. Kobayashi and T. Maskawa, Prog. Theor. Phys. 49(2), 652 (1973); N. Cabibbo, Phys. Rev. Lett. 10, 531 (1963).

[3] M. Gronau and D. Wyler, Phys. Lett. B 265, 172 (1991); M. Gronau and D. London, Phys. Lett. B 253, 483 (1991).

[4] D. Atwood, I. Dunietz and A. Soni, Phys. Rev. D 63, 036005 (2001); D. Atwood, I. Dunietz and A. Soni, Phys. Rev. Lett. 78, 3257 (1997).

[5] A. Giri, Y. Grossman, A. Soffer, and J. Zupan, Phys. Rev. D 68, 054018 (2003).

[6] R. Fleischer, Phys. Lett. B 459, 306 (1999); R. Fleischer and J. Matias, Phys. Rev. D 66, 054009 (2002).

[7] I. Dunietz, R. Fleischer and U. Nierste, Phys. Rev. D 63, 114015 (2001).

[8] E. Barberio et al. (Heavy Flavor Averaging Group (HFAG)), arXiv:hep-ex/0808.1297v3 (2007).

[9] B. Aubert et al. (BABAR Collaboration), Phys. Rev. D 77, 111102 (2008).

[10] J. Charles et al. (CKMfitter Group) Eur. Phys. J. C 41, 1 (2005), updated results and plots available at: http://ckmfitter.in2p3.fr; M. Bona et al. (UT-fit Collaboration) JHEP, 0610:081 (2006), arXiv:hep-ph/0601677v2, updated results and plots available at: http://www.utfit.org/.

[11] Y. Grossman, A. Soffer and J. Zupan, Phys. Rev. D 72, 031501(R) (2005).

[12] M. Gronau, Phys. Lett. B 557, 198 (2003); M. Gronau, Phys. Rev. D 58, 037301 (1998).

[13] D. E. Acosta et al. (CDF Collaboration), Phys. Rev. D 71, 032001 (2005).

[14] A. Sill, Nucl. Instrum. Methods A 447, 1 (2000).

[15] A. Affolder et al., Nucl. Instrum. Methods A 518, 532 (2004).

[16] CDF II uses a cylindrical coordinate system in which \( \phi \) is the azimuthal angle, \( r \) is the radius from the nominal beam line, and \( z \) points in the proton beam direction, with the origin at the center of the detector. The transverse plane is the plane perpendicular to the \( z \) axis.

[17] E. J. Thomson et al., IEEE Trans. Nucl. Sci. 49, 1063 (2002).

[18] B. Ashmanskas et al., Nucl. Instrum. Methods A 518, 532 (2004).

[19] Isolation is defined as \( I_B = p_T(B)/(p_T(B) + \sum p_T_i) \), where \( p_T(B) \) is the transverse momentum of the \( B \) candidate, and the sum runs over all other tracks within a cone of radius 1, in \( \eta - \phi \) space around the \( B \) flight-direction. Its value is typically higher for bottom–flavored hadrons than for random track combinations.

[20] A. Abulencia et al. (CDF Collaboration), Phys. Rev. Lett. 97, 211802(2006).

[21] T. Aaltonen et al. (CDF Collaboration), Phys. Rev. Lett. 103, 031801(2009).

[22] P. Squillacioti, Ph.D. Thesis, University of Siena, FERMILAB-THESIS-2006-27, Nov 2006.

[23] R. Brun, R. Hagelberg, M. Hansroul, and J. C. Lasalle, Reports no. CERN-DD-78-2-REV and CERN-DD-78-2.

[24] D. E. Acosta et al. (CDF Collaboration), Phys. Rev. Lett. 94, 122001 (2005).

[25] M. Morello (CDF Collaboration), Nucl. Phys. Proc. Suppl. 170, 39 (2007).

[26] D. E. Acosta et al. (CDF Collaboration), Phys. Rev. Lett. 96, 202001 (2006).