Independence and Product Systems

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Abstract. Starting from elementary considerations about independence and Markov processes in classical probability we arrive at the new concept of conditional monotone independence (or operator-valued monotone independence). With the help of product systems of Hilbert modules we show that monotone conditional independence arises naturally in dilation theory.

One of the most fundamental concepts in probability theory — in the sequel, we say classical probability — is independence. The most fundamental concept dealing with non-independent random variables is the Markov property. In order to underline the formal aspect of the Markov property being a generalization of independence (replacing expectations by conditional expectations) we will say conditional independence (which operator algebraists would prefer to call operator-valued independence).

The expectation is a normalized positive linear functional (i.e. a state) on the commutative algebra of random variables. Noncommutative or quantum probability is a noncommutative generalization of classical probability designed to include also quantum physical applications. A quantum probability space is, therefore, a pair (A, φ) of a unital *-algebra (more specifically, a (pre-)*C*-algebra or a von Neumann algebra) A with a normalized (i.e. φ(1) = 1) positive (i.e. φ(a* a) ≥ 0) linear functional φ: A → C. In this context a classical probability space (Ω, F, P) corresponds to the commutative quantum probability space (L∞(Ω), φ = ∫ φ(dP)).

Classical independence is symmetric, i.e. if X₁ is independent of X₂ (conditioning X₁ on X₂ does not change probabilities of X₁), then X₂ is independent of X₁, too. In Section 1 we discuss an example where this symmetry is not desirable. However, within the category of commutative quantum probability spaces there is only one independence, namely, that one which gives back classical independence. Therefore, if we want to model the mentioned example, then we are forced to leave the category of commutative quantum probability spaces even for classical random variables. The only noncommutative independence which earns the name independence and is not symmetric is monotone independence; see Section 1.

In Section 2 we discuss the case of conditional independence, which is the basis for our discussion of Markov processes, later on. As a new concept we introduce conditional monotone independence (or operator-valued monotone independence). As long as we are dealing with classical random variables, both conditional independence and conditional monotone independence are possible. However, if we
start with noncommutative random variables from the beginning, then the constructions for conditional independence are no longer possible, in general. Only in few exceptional cases (for instance, when the involved quantum probability spaces are algebras isomorphic to some algebra $\mathcal{B}(H)$ of all bounded operators on some Hilbert space $H$) quantum conditional independence still exists.

In Section 3, finally, we point out with the help of product systems of Hilbert modules that conditional monotone independence arises in every reversible quantum dynamical system which is a dilation of an irreversible quantum dynamical system. (Conditional) monotone independence is, therefore, one of the most natural quantum independences possible.

1. Independence

Let $X_1, X_2$ stand for two observable quantities. These can be classical random variables (i.e. real-valued measurable functions on some probability space) or quantum random variables (i.e. self-adjoint operators on some Hilbert space with a distinguished vector state $\varphi = \langle \xi, \cdot \xi \rangle$).

To avoid technical discussions, we assume that the $X_i$ are bounded. (This is not a serious restriction, because on the classical side everything is determined when known on indicator functions, and on the quantum side by the spectral theorem everything is known when known on projections.)

Let us start with the case of two classical random variables. Everything one can say in the classical description can be formulated in terms of (marginal and joint) distributions.

$$X_i \leadsto L^\infty(\mathbb{R}, \mu_i) \subset L^1(\mathbb{R}, \mu_i)$$ marginal distribution

$$\mathbb{E}[f(X_i)] = \int f(x) \mu_i(dx)$$

$$X_1, X_2 \leadsto L^\infty(\mathbb{R} \times \mathbb{R}, \mu_{12})$$ joint distribution

$$\mathbb{E}[f(X_1, X_2)] = \int f(x_1, x_2) \mu_{12}(dx_1, dx_2)$$

The quantum description of the classical case is the following reinterpretation.

$$\mathcal{A}_i = L^\infty(\mathbb{R}, \mu_i) \subset \mathcal{B}(L^2(\mathbb{R}, \mu_i))$$ algebras

$$\varphi_i(f) = \int f(x) \mu_i(dx) = \langle 1, f1 \rangle$$ with a state

$$\mathcal{A}_{12} = L^\infty(\mathbb{R} \times \mathbb{R}, \mu_{12})$$ a “product” of algebras

$$\varphi_{12}(f) = \int f(x_1, x_2) \mu_{12}(dx_1, dx_2) = \langle 1, f1 \rangle$$ with a state on the “product”

$X_1$ and $X_2$ are independent, if $\mu_{12} = \mu_1 \otimes \mu_2$. Then

$$\mathcal{A}_{12} = \mathcal{A}_1 \otimes \mathcal{A}_2 \quad \varphi_{12} = \varphi_1 \otimes \varphi_2$$

$$L^2(\mathbb{R} \times \mathbb{R}, \mu_{12}) = L^2(\mathbb{R}, \mu_1) \otimes L^2(\mathbb{R}, \mu_2)$$

$$1 = 1 \otimes 1.$$
simultaneously in such a way that they influence each other as little as possible. In the preceding discussion, independence, formulated intrinsically as factorization of probabilities, gives rise to a tensor product. The tensor product, in turn, can also serve a prescription how to construct a joint probability space from the marginal distributions.

Let us start with the constructive interpretation. Schürmann [Sch95] shows that the tensor product of unital commutative quantum probability spaces \((A_i, \varphi_i)\) is the only prescription, how to compute mixed moments \(\varphi_{12}(X_1^n X_2^m)\), such that certain axioms are fulfilled. These axioms are properties derived from classical independence like associativity of the construction and the requirement that functions of independent random variables should again be independent (naturality).

We see that independence inside the category of commutative quantum probability spaces is fixed uniquely by natural properties which nobody wants to miss. However, coming back to the intrinsic interpretation, although the statement “X2 does not influence X1” is not symmetric under exchange of X1 and X2, classical independence is symmetric. In other words, within a single classical probability space there is no possibility to find random variables X1, X2 such that X2 is independent of X1 but not conversely. Is this realistic? In other words, are all situations in which we want to derive probabilistic results of the type that, if measurement of X2 does not influence knowledge on X1, then measurement of X1 does not influence knowledge on X2?

We consider an example. Suppose X1 is an observable at instant \(t_1\) and X2 is an observable at instant \(t_2 > t_1\). We naturally expect that X1 is independent of X2: (A measurement at time \(t_2\) cannot influence earlier times.) X2 may depend on whether we measured X1 or not. In the extreme case, measuring X1 “destroys” what we know about X2. (Conditioning X2 on X1, we obtain the constant function \(E(X_2)\) of X1.)

Can we model this situation by a classical model (i.e. commutative algebras)? No, certainly not, because classical independence is symmetric. Can we model the situation by a noncommutative model (Xi represented as operators on a Hilbert space, say)? Yes, as follows. As before, we represent X2 by the operator

\[
\text{id} \otimes X_2 \text{ on } L^2(\mathbb{R}, \mu_1) \otimes L^2(\mathbb{R}, \mu_2).
\]

For X1, however, we choose

\[
X_1 \otimes 1^* \ (1^* = |1\rangle\langle 1| : f \mapsto 1(1, f)).
\]

We find

\[
(1.1) \quad \varphi_{12}(g_0(X_2) f_1(X_1) g_1(X_2) \ldots f_n(X_1) g_n(X_2)) = \mathbb{E}[g_0(X_2)] \mathbb{E}[g_1(X_2)] \ldots \mathbb{E}[g_n(X_2)] \cdot \mathbb{E}[(f_1 \ldots f_n)(X_1)].
\]

In other words, the (quantum) random variables X1 and X2 (more precisely, the algebras \(A_1 \otimes 1^*\) and \(\text{id} \otimes A_2\) generated by them) are monotone independent in the sense of Lu [Lu97] and Muraki [Mur97].

A possible interpretation: Given \(f(X_1)\), measuring \(g(X_2)\), the information on X1 is fully preserved. Indeed, as the expression

\[
f'(X_1) g(X_2) f(X_1) = (f' f)(X_1) \mathbb{E}[g(X_2)]
\]

shows, a second measurement \(f'(X_1)\) “sees” the full function \(f(X_1)\) so that the product \((f' f)(X_1)\) occurs. On the other hand, the measurement of \(f'(X_1)\) has completely “nivelated” any information about X2 present after the measurement of
Provocantly, for the conditioning on \( X_1 \) under the condition \( X_2 \in B \) we might write
\[
E(X_1|X_2 \in B, X_1) = X_1 \ E(X_2|X_2 \in B, X_1) = 1.
\]
We see that a very natural model involving only classical random variables \( X_1, X_2 \) leads naturally to a (noncommutative) quantum probability space.

The commutative nature of \( A_i \) does not play a role. Allowing for noncommutative algebras we obtain immediately (quantum) tensor independence and (quantum) monotone independence. (Observe that the representation space for the joint distribution is the same in both cases. Only the embedding of \( A_i \) — unital in the tensor case, and non-unital in the monotone case — differs.)

There are many other (quantum) independences, in the first place, Voiculescu’s famous free independence and boolean independence (von Waldenfels [Wal73]). Additionally, many authors consider notions of convolutions which lead to (quantum) central limit theorems as independences. However, by the following theorem we see that only the four mentioned independences fulfill Schürrmann’s axioms. (\( q \)-Convolution, for instance, is not natural; see van Leeuven and Maassen [LM95].)

1.1 Theorem (Speicher [Spe97], Ben Ghorbal and Schürrmann [BGS99], Muraki [Mur02]). There are only two unital and symmetric independences, namely, tensor and free independence, fulfilling Schürrmann’s axioms. There is only one (family of) non-unital, symmetric independence(s), namely boolean independence. There is only one non-unital, non-symmetric independence, namely, monotone independence.

1.2 Remark. In our example we decided that measuring \( X_1 \) nivelates information about \( X_2 \) while measuring \( X_2 \) leaves unaffected information about \( X_1 \). Of course, there are many thinkable ways how \( X_2 \) can influence \( X_1 \). However, the preceding theorem tells us that only the nivelation as described in our example can be interpreted as (quantum) independence.

2. Conditional independence

Let us consider a (classical) game: Two players 1,2 throw three coins, one “fair” coin (\( p_{\text{head}} = p_{\text{tail}} \)), one “good” coin (favouring head) and one “bad” (favouring tail), in the following way. First, they throw the fair coin. Depending on the result, player 1 throws the “good” coin and player 2 throws the “bad” one, or conversely. Denote by

\[
\begin{align*}
Y & \quad \text{the outcome of the first throw (of the “fair” coin),} \\
X_i & \quad \text{the outcome of the second throw for player } i.
\end{align*}
\]

Of course, \( X_1 \) and \( X_2 \) are not independent. They are, however, independent given that \( Y \) has a certain value. In other words, they are conditionally independent (over \( Y \)), i.e.

\[
E(f(X_1)g(X_2)|Y) = E(f(X_1)|Y) \ E(g(X_2)|Y).
\]

Equivalently, the conditional probabilities factorize

\[
P(y, dx_1 \times dx_2) = P(y, dx_1) \ P(y, dx_2).
\]

We make two observations:
• Given the probability spaces \((\mathbb{R} \times \mathbb{R}, \mu_0)\) describing each pair \((Y, X_i)\) (or, equivalently, the corresponding \(L^2\)-spaces) it seems difficult to construct (by an abstract construction involving only the abstract structure of probability spaces) the joint probability space \((\mathbb{R} \times \mathbb{R} \times \mathbb{R}, \mu)\) for \((Y, X_1, X_2)\) directly from the former.

• The above factorizations do not depend on the actual distribution \(\nu\) of \(Y\), but only on its \(\sigma\)-algebra, i.e. on the algebra \(L^\infty(\mathbb{R}, \nu)\).

Quantum probability can help solving the problem in the first issue taking into account the second. What do we have? We have two pairs \((A_i, \Phi_i)\)

\[
A_i = L^\infty(\mathbb{R} \times \mathbb{R}, \mu_0) \quad \text{an algebra}
\]

\[
\Phi_i: A_i \to A_0 = L^\infty(\mathbb{R}, \nu) \subset A_i \quad \text{a conditional expectation}
\]

with \(\Phi_i(f(Y)g(X_i)) = \mathbb{E}(f(Y)g(X_i)|Y) = \mathbb{E}(g(X_i)|Y)f(Y)\). Our task is to construct a joint (quantum) probability space with a conditional expectation \(\Phi\) factorizing according to \(\mathbf{2.1}\). The problem: Be it the product of classical probability spaces or be it the tensor product of quantum probability spaces, in both products we end up with two copies of \(Y\) (or better of the subalgebra \(A_0\)) contained in each factor.

How can we identify suitably members of the two different copies? The answer follows from the observation that by the bimodule property of conditional expectations, the factorization \(\mathbf{2.1}\) does not “see” from which factor \(A_i\) an element \(h(Y)\) comes:

\[
\mathbb{E}(f(X_1)h(Y)g(X_2)|Y) = \mathbb{E}(f(X_1)h(Y)|Y) \mathbb{E}(g(X_2)|Y) = \mathbb{E}(f(X_1)|Y) \mathbb{E}(h(Y)g(X_2)|Y) = \mathbb{E}(f(X_1)|Y) \mathbb{E}(h(Y)g(X_2)|Y).
\]

Therefore, we can amalgamate over \(A_0\) and consider the \(A_0-A_0\)-module tensor product of the \(A_0-A_0\)-modules \(A_i\), instead of the usual tensor product. (As mentioned this construction depends only on the algebra \(A_0\) and the conditional expectations \(\Phi_i\), but not on \(\nu\).) Then the module tensor product

\[
A_1 \circ A_2 := A_1 \otimes A_2 / a_1a_0 \otimes a_2 - a_1 \otimes a_0a_2
\]

is an algebra (with multiplication \((a_1 \otimes a_2)(a'_1 \otimes a'_2) = a_1a'_1 \otimes a_2a'_2\)) isomorphic to (a dense subalgebra of) \(L^\infty(\mathbb{R} \times \mathbb{R} \times \mathbb{R}, \mu)\) and \(\Phi(a_1 \otimes a_2) = \Phi_1(a_1)\Phi_2(a_2)\) defines a conditional expectation such that

\[
\Phi(f(X_1)g(X_2)) = \mathbb{E}(f(X_1)|Y) \mathbb{E}(g(X_2)|Y).
\]

In other words, \(A_1\) and \(A_2\) are tensor independent \textbf{with amalgamation} (over \(A_0\)) in the sense of Skeide \cite{Ske93a, Ske99}. The terminology “with amalgamation” follows \cite{Voi95, Spe98}. Today we would prefer to say \(A_1\) and \(A_2\) are \textbf{conditionally tensor independent}.

On the level of representation spaces the product of probability spaces was reflected by the tensor product. We now explain that for conditional independence the “good” representation spaces are \textbf{Hilbert} \(A_0-\mathbb{A}_0\)-\textbf{modules} and the “correct"
product construction, once again, is the tensor product over $A_0$. Let

$$E_i = \frac{A_i}{N_i}$$

with inner product

$$\langle x, x' \rangle = \Phi_i(x^* x')$$

so that

$$\Phi_i(a) = \langle 1, a1 \rangle$$

where $N_i$ indicates the submodule of length-zero elements. (This is Paschke’s [Pas73] GNS-construction for CP-maps applied to the simpler case of conditional expectations as in Rieffel [Rie74].) The tensor product over $A_0$

$$E_1 \otimes E_2 = \frac{E_1 \otimes E_2}{N}$$

(2.2) with inner product

$$\langle x_1 \otimes x_2, x_1' \otimes x_2' \rangle = \langle x_2, \langle x_1, x_1' \rangle x_2' \rangle$$

carries a natural representation of $A_1 \otimes A_2$ (inherited from $A_1 \otimes A_2$) and we recover $\Phi$ as $\Phi(a) = \langle 1, a1 \rangle$ (where $1 = 1 \otimes 1$).

2.1 Remark. Of course, for our special case in (2.2) we might write also $\langle x_1 \otimes x_2, x_1' \otimes x_2' \rangle = \langle x_2, \langle x_1, x_1' \rangle x_2' \rangle$ for the inner product, because everything commutes and $\Phi_2$ is a conditional expectation. However, if we write the definition of the inner product of a tensor product as in (2.2), then it works for arbitrary Hilbert bimodules. It seems to be a general feature that writing things in the “correct” order helps both generalizing to the noncommutative case and also understanding the classical case better.

What about conditional tensor independence for arbitrary quantum probability spaces? Here we run into serious problems. Given two algebras $A_i$ with a subalgebra $A_0$, there is only in exceptional cases a reasonable multiplication on $A_1 \otimes A_2$. (For instance, if $A_0 = \mathcal{B}(H)$.) Likewise, in general,

$$E_1 \otimes E_2 \not\cong E_2 \otimes E_1 \quad \exists A_1 \longrightarrow A_1 \otimes \text{id} \quad \# A_2 \longrightarrow \text{id} \circ A_2.$$

This shows that, in general, conditional tensor independence does not have an analogue in quantum probability.

2.2 Remark. The non-isomorphism $E_1 \otimes E_2 \not\cong E_2 \otimes E_1$ is at the heart of the non-possibility to construct the amalgamated analogue of tensor independence. For instance, if one of the tensor products is the null space, then certain moments, computed according to the rules in the proof of Theorem 1, must be $= 0$ and $\neq 0$ simultaneously. At least, in the case, of von Neumann algebras plus the condition $E_1 \otimes E_2 \cong E_2 \otimes E_1$, we think that it is possible to overcome the difficulties and to define the analogue of tensor independence. Notice, however, that there is no flip $x_1 \otimes x_2 \mapsto x_2 \otimes x_1$ available. The embedding $A_2 \hookrightarrow \mathcal{B}(E_1 \otimes E_2)$ is not that commuting from the isomorphism $E_1 \otimes E_2 \to E_2 \otimes E_1$ but it is considerably more subtle.

However, there exists always an embedding $A_2 \to "11^* \otimes A_2"$, where "11* $\otimes A_2$" is not a completely exact abbreviation for operators of the form

(2.3) $$(11^* \otimes a_2)(x_1 \otimes x_2) := (1 \otimes \text{id})a_2(1^* \otimes \text{id})(x_1 \otimes x_2) = 1 \otimes a_2(1, x_1)x_2.$$ (It is clear that in the scalar case $A_0 = \mathbb{C}$ this reduces to the well-defined operator $11^* \otimes a_2$ from Section 1).

Since $(11^* \otimes a_2)(a_1 \otimes \text{id})(11^* \otimes a_2') = 11^* \otimes a_2\Phi(a_1)a_2'$, we find

(2.4) $$\Phi(a_1^{(0)} a_2^{(1)} a_1^{(1)} \ldots a_2^{(n)} a_1^{(n)}) = \Phi(a_1^{(0)}) \Phi(a_2^{(1)} \Phi(a_1^{(1)} \ldots a_2^{(n)}) \Phi(a_1^{(n)}).$$
By analogy with (1.1) we propose a new definition such that the algebras $A_1 \odot \text{id}$ and $11^* \odot A_2$ become conditionally monotone independent over $A_0$.

**2.3 Definition.** Let $A$ be a $*$-algebra with a conditional expectation $\Phi$ onto the $*$-subalgebra $A_0$. An ordered pair $(A_1, A_2)$ of $*$-subalgebras $A_1, A_2$ containing $A_0$ is **conditionally monotone independent** (over $A_0$ in $\Phi$), if they fulfill (2.4).

**2.4 Remark.** Comparing with our scalar example of monotone independence, we see that $A_1$ and $A_2$ have changed their roles. The non-unitaly embedded algebra (standing for events in the past of the other) is now $A_2$. Contrary to all good habits, the module picture shows:

> The future is on the left of the past!

We close the discussion of conditional independence with the analogue of Theorem 1. Every independence except tensor independence has an analogue conditional independence. Therefore, in Theorem 1 we have just to cancel tensor independence (of course, there cannot be more conditional independences, because this is not possible already in the scalar case) and arrive at the pleasant situation where every type of independence has now exactly one realization.

**2.5 Theorem.** There is exactly one **unital** conditional independence, namely, conditional **free** independence (which is **symmetric**, automatically). There is exactly one **non-unital, symmetric** conditional independence, namely conditional **boolean** independence. There is exactly one **non-symmetric** conditional independence, namely, conditional **monotone** independence (which is **non-unital**, automatically).

Conditional free independence was introduced by Voiculescu [Voi95] and studied intensively by Speicher [Spe98]. Conditional boolean independence is mentioned in Skeide [Ske00] (without the simple proof of existence).

### 3. Dilations

Reversible and irreversible (quantum) dynamical systems are related by dilations. A dilation of a (small) irreversible system is a (big) reversible system such that the irreversible dynamics of the small system can be understood as a projection (conditional expecation) from the reversible dynamics of the big system onto the small subsystem. We illustrate this.

Small system: A unital $C^*$-algebra $B$

Irreversible evolution: A CP-semigroup $T$

Big system: A unital $C^*$-algebra $A$

Reversible evolution: An $E_0$-semigroup $\varphi$

A CP-semigroup is a semigroup of completely positive mappings, while an $E_0$-semigroup is a semigroup of unital endomorphisms.

Both the upper and the lower half of the diagram are closely related to and, actually, connected by tensor product systems of Hilbert modules as introduced in
Bhat and Skeide [BS00]. We explain this in the case of a classical Markov process (joint work with L. Accardi to be published elsewhere).

We consider

\[ \mathcal{B} = L^\infty(\mathbb{R}, \nu) \]

\[ [T_t(f)](y) = \int P_t(y, dx) f(x) \]

\[ \mathcal{A} = L^\infty(\prod_{t \in \mathbb{R}_+} \mathbb{R}, \mu) \]

\[ \vartheta_t(X_s) = X_{s+t} \]

\[ \iota(Y) = X_0 \]

\[ p = \mathbb{E}(\bullet ; Y). \]

Once again, what we do in the sequel does not depend on the measure \( \nu \) but only on its measure type. The technical condition of \( P_t(x, \bullet) \) being absolutely continuous with respect to \( \nu \) assures that \( T_t \) is well-defined. In practise, one will even choose \( \nu \) according to this requirement, even if the distribution of \( X_0 \) will be different. (For instance, it could be \( \delta_0 \); see Remark 3.)

We define a filtration. For \( I \subset \mathbb{R}_+ \) set \( \mathcal{A}_t = L^\infty(\prod_{t \in I} \mathbb{R}, \mu_I) \) and \( \mathcal{A}_t = \mathcal{A}_{[0,t]} \).

\[ E = \mathcal{A}/\mathbb{N} \]

\[ E_t = \mathcal{A}_t/\mathbb{N}_t \]

\[ \langle x, x' \rangle = p(x^*x') \]

Then \( E \) is a Hilbert \( \mathcal{B} \)-module, \( E_t \) are Hilbert \( \mathcal{B} \)-modules (left action via \( X_t \), i.e. \( f \in L^\infty(\mathbb{R}, \nu) \) acts as multiplication by the function \( f(X_t) \)) such that

\[ E \otimes E_t = E \]

\[ E_t \otimes E_t = E_{s+t} \]

\[ \vartheta_t(a) = a \otimes id_{E_t} \]

\[ 1 \otimes 1_t = 1 \]

\[ 1_s \otimes 1_t = 1_{s+t} \]

\[ T_t(b) = (1, \vartheta_t \circ \iota(b)) = (1_t, b_1) \]

For instance, looking at the definition of the tensor product in [22], one checks that

\[ (f_n(X_{p_n}) \ldots f_1(X_{p_1})) \otimes (g_m(X_{q_m}) \ldots g_1(X_{q_1})) \]

\[ \longleftrightarrow f_n(X_{p_n} + t) \ldots f_1(X_{p_1 + t}) g_m(X_{q_m}) \ldots g_1(X_{q_1}) \]

defines an isomorphism \( E \otimes E_t \to E \). The restrictions to \( E_s \otimes E_t \) maps onto \( E_{s+t} \subset E \) and also the left multiplication is preserved.

In particular, \( E^\circ = (E_t) \) is a product system \((E_s \otimes E_t = E_{s+t}, E_0 = \mathcal{B}) \) and \( 1^\circ = (1_t \otimes 1_t = 1_{s+t}, l_0 = 1) \) for \( E^\circ \).

Of course, the \( E_0 \)-semigroup \( \vartheta \) extends to all of \( \mathcal{B}^a(E) \) (still sending \( a \) to \( a \otimes id_{E_t} \)). In Skeide [Ske02] we showed that for every \( E_0 \)-semigroup on some \( \mathcal{B}^a(E) \) (\( \vartheta \) being continuous in a certain topology called strict) we can construct a (unique) product system \( E^\circ \) such that \( \vartheta_t \) can be recovered from a factorization \( E = E \otimes E_t \).

(Actually, we require existence of a unit vector \( \xi \in E \), i.e. \( \langle \xi, \xi \rangle = 1 \). The general case can be treated for von Neumann algebras; see [MSS03]. The construction in [Ske02] generalizes a construction for Hilbert spaces from Bhat [Bha96], while the construction in [MSS03] generalizes the construction of product systems of Hilbert spaces from Arveson [Arv89].)

We return to the construction of product systems from CP-semigroups on general \( C^* \)-algebras, and we want to see how monotone independence arises naturally.
In [BS00] we had to face a problem, namely, the possibility of a unital embedding \( B \to \mathcal{A} = \mathcal{B}^a(E) \) is restricted to commutative \( B \). Instead, we had to use \( i(b) = \overline{b}^{\ast} \) leading to a different filtration. Setting \( \hat{\mathcal{A}}_t = \text{alg}(\vartheta_s(\mathcal{A}_t)) \) \((0 \leq s \leq t)\). While \( \mathcal{A}_t \) is commutative and, therefore, cannot coincide with \( \mathcal{B}^a(E_t) \), we find that (a sufficient closure of) \( \hat{\mathcal{A}}_t \) is isomorphic to \( \mathcal{B}^a(E_t) \) embedded as \( \hat{\mathcal{A}}_t = 1^* \odot \mathcal{B}^a(E_t) \). Also the closure of \( \hat{\mathcal{A}}_\infty \) is all of \( \mathcal{B}^a(E) \).

Neither the algebras \( \mathcal{A}_t \) and \( \hat{\mathcal{A}}_{[t,t+s]} \) are conditionally tensor independent nor the algebras \( \hat{\mathcal{A}}_t \) and \( \hat{\mathcal{A}}_{[t,t+s]} = \vartheta_t(\hat{\mathcal{A}}_s) \) are conditionally monotone independent. For that the dilation must be a white noise, i.e. a dilation of the trivial semigroup \( T_t = \text{id}_B \), in yet other words, the conditional expectation \( p = (1, \cdot 1) \) must be invariant for \( \vartheta \). The same computations leading to (2.4) show also the following.

3.1 Theorem. Let \((E, \vartheta, \xi)\) be a white noise, i.e. \( \vartheta \) is a (strict) \( E_0 \)-semigroup on \( \mathcal{B}^a(E) \) leaving \( p = (1, \cdot 1) \) invariant. Set \( \hat{\mathcal{A}} = \xi \xi^* \odot \mathcal{B}^a(E_t) \) (cf. (2.3)) and \( \hat{\mathcal{A}}_{[t,t+s]} = \vartheta_t(\hat{\mathcal{A}}_s) \). Then for all \( 0 < r < s < t \) the algebras \( \hat{\mathcal{A}}_{[r,s]} \) and \( \hat{\mathcal{A}}_{[s,t]} \) are conditionally monotone independent over \( \hat{\mathcal{A}}_0 \).

3.2 Remark. Even if \( \vartheta \) is not a white noise, then \( \hat{\mathcal{A}}_{[r,s]} \) and \( \hat{\mathcal{A}}_{[s,t]} \) are conditionally monotone independent over \( \hat{\mathcal{A}}_s \). This is just the (quantum) Markov property.

3.3 Remark. If \( T_t = \text{id}_B \) then \( E^\otimes \) constructed as in [BS00] is trivial \( (E_t = B) \). However, most dilations are cocycle perturbations of a (non-trivial, if \( T \) is non-trivial) white noise, and cocycle perturbation does not change the product system; see [Ske02].

If there is a (non-trivial) white noise, then by Theorem 8 there is always a (non-trivial) filtration of subalgebras \( \mathcal{A}_t \) such that the increment algebras \( \mathcal{A}_{s,t} \) to disjoint intervals are conditionally monotone independent. This justifies, in particular, the name white noise.

3.4 Remark. In the scalar case \( B = \mathbb{C} \) all (unital) CP-semigroups are trivial. Therefore, if \( \vartheta \) is an \( E_0 \)-semigroup on some \( \mathcal{B}(\mathcal{H}) \) (Arveson [Arv89]) and \( \xi \in \mathcal{H} \) a unit vector such that \( \vartheta_t(\xi^* \xi) \) is increasing, then the increment algebras \( \vartheta_t(\text{id}_B \odot \mathcal{B}(\mathcal{H}_t)) \) and \( \vartheta_t(\omega \omega^* \odot \mathcal{B}(\mathcal{H}_t)) \) are tensor independent and monotone independent, respectively.

3.5 Remark. Suppose \( X = (X_t) \) is a stationary independent increment process. Thus, \( Y = (Y_t) \) with \( Y_t = X_t - X_0 \) is a Lévy process. Then \( Y_0 \sim \delta_0 \). Since \( L^\infty = \bigotimes (\mathbb{R}, \delta_0) = \mathbb{C} \), the associated product system consists of Hilbert spaces. It is not difficult to see that the \( L^2 \)-space of \( Y \) is \( \mathcal{H} = \Gamma(L^2([0,t], K)) \) and the associated product system is \( \mathcal{H}_t = \Gamma(L^2([0,t], K)) \) (Parthasarathy and Schmidt [PS72]). The product system \( E^\odot \) of \( X_t \) is intimately but non-trivially related to \( \mathcal{H} \) (joint work with U. Franz and M. Schürmann; unpublished).

For a more detailed survey on product systems of Hilbert modules and the present status of their classification we refer the reader to Skeide [Ske03b].

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