Dark matter signals and cosmic ray anomalies in an extended seesaw model

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\textbf{A B S T R A C T}

An extended seesaw model proposed to achieve low scale leptogenesis can resolve the excess positron and electron fluxes observed from PAMELA, ATIC and/or Fermi-LAT, and simultaneously accommodate some of recent experimental results for dark matter (DM) signals. In this approach, in addition to $SU(2)_L$ doublet and the (light) singlet Higgs fields, an extra vector-like singlet neutrino and a singlet scalar field, which are coexisting two-particle dark matter candidates, are responsible for the origin of the excess positron and electron fluxes to resolve the PAMELA, ATIC and/or Fermi-LAT anomalies, as well as for the DM signals observed from direct searches in low mass scale.

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1. Introduction

The quest for identification of the missing mass of our universe is one of the most fundamental issues in astroparticle physics and cosmology. The evidence for non-baryonic dark matter (DM) inferred from a combination of cosmological and astrophysical phenomena becomes more and more convincing, which alludes the existence of new physics beyond the standard model (SM). Very recently, several new exciting data on DM have been released, which may open up new era to search for DM in a low mass region of a few GeV. CDMS II Collaboration reported the two DM candidate events with a 77% C.L and the upper bound of null result [1]. DAMA Collaboration confirmed the model independent evidence of the presence of DM on the basis of the DM annual modulation signature with 8.9$\sigma$ significance [2]. The CoGeNT experiment reported a possible signal of a light DM candidate with $m_{\text{DM}}=7$–11 GeV, and provided 90$\%$ C.L. WIMP exclusion plots as well [3]. Those three independent experimental results may be interpreted as signals of the existence of DM with a low mass around a few GeV [4]. Contrary to the results from CDMS II, DAMA and CoGeNT, XENON100 Collaboration announced that they have not observed any DM signal for the similar parameter ranges searched by those three experiments [5]. Therefore, we need further experimental results to judge if there really exists a DM candidate with a low mass or not.

On the other hand, the PAMELA experiment has presented a significant positron flux excess over the expected background with no excess in the corresponding anti-proton flux [6]. The ATIC/PPB-BETS experiment has shown significant excess of electron and positron flux at energies around 300–800 GeV [7,8]. More recently, Fermi-LAT experiment have also shown an excessive electron and positron flux in the same energy range as in ATIC but its strength was not strong compared to ATIC [9]. So, it is likely that the experimental evidences for the signals of DM with a low mass scale are not reconciled with the cosmic ray positron and electron excess in the framework of one and only one DM scenarios.

Recently, we have proposed an extended seesaw model to simultaneously and naturally accommodate tiny neutrino masses, low scale leptogenesis and dark matter candidate by introducing extra singlet neutrinos and singlet scalar particles on top of the canonical seesaw model [10,11]. Furthermore, we have proposed a coexisting two-particle DM scenario [12] by allowing both an extra singlet Majorana neutrino and a light singlet scalar particle as two DM candidates. Such a scenario containing more than one DM may be desirable in the case that there exist a few incompatible phenomena which are very hard to reconcile in the scenarios with only one DM.

The purpose of this Letter is to investigate how both the low mass DM signals observed from direct DM searches and the cosmic ray positron and electron excess observed from PAMELA, ATIC and/or Fermi-LAT experiments are simultaneously explained in the extended seesaw model with coexisting two-particle DM proposed in [12]. Due to the tension among the experimental results of direct search for DM in low mass scale, we first consider the case that lighter DM candidate in our model has mass around 3 GeV.
allowed by DAMA experiment, which is not in conflict with other null results from direct searches but is inconsistent with the DM signals observed from CoGeNT. The other case we consider is to accept DAMA and CoGeNT signals for DM candidate whose overlapped mass range lies between 7 GeV and 11 GeV while ignoring XENON100 results. In this work, we slightly modify the model proposed in [12] by replacing extra singlet Majorana neutrino with singlet vector-like neutrinos so as to simply resolve the cosmic ray anomaly while keeping to accommodate tiny neutrino masses and low scale leptogenesis of order 1–10 TeV [10,11].

We notice that to achieve our coexisting two-particle DM scenario in the renormalizable framework as shown in [12], an extra singlet Higgs scalar field $\Phi$ is necessarily introduced, which may open up new channels of DM annihilations. As will be shown later, in this scenario, this scalar field $\Phi$ may play an essential role in resolving the unexpected electron and positron fluxes measured at PAMELA, ATIC and/or Fermi-LAT if the mass of $\Phi$ has rather small around just below 1 GeV so as for the annihilation cross section to be enhanced via a mechanism first described by Sommerfeld [13–15]. Once this new force carrier $\Phi$ is included, the possibility of a new dominant annihilation of singlet vector-like neutrinos into a pair of $\Phi$ opens up. The $\Phi$ mixes with the Higgs allowing it to decay into the final state fermions, and if the $\Phi$ is taken to be light, it is kinematically constrained to decay to mostly lepton pairs preventing from producing anti-protons, so that the excess of positron and/or electron observed can be accounted for. In addition, the low mass DM signals will be explained by considering the singlet scalar $\psi$ as the lightest DM candidate with mass of order a few GeV. Thus, the low mass DM signals, the excess positron and electron fluxes produced from the cosmic rays, low scale leptogenesis and light neutrino masses can be simultaneously accommodated in our model proposed.

To see how the coexisting two-particle DM scenario is achieved, let us consider the following Lagrangian

$$
L = L_0 + (Y_D L H N + Y_\gamma N \psi S + \text{h.c.}) + M_N N^T N + Y_\Phi S \Phi S - m_\phi S \Phi S + \frac{1}{2} m_\phi \Phi^2 - \frac{\lambda_\phi}{2} \Phi^4 - m_\lambda \lambda^2 \Phi^2 - m_\lambda H^\dagger H \Phi^2,
$$

(1)

where the first term is the Lagrangian of the SM and kinetic terms of the singlet fields, and $L$, $N$, $S$ and $\psi$ stand for $SU(2)_L$ lepton doublet, singlet heavy Majorana neutrino, singlet vector-like neutrino composed of two Weyl fermions, and light singlet scalar, respectively. Note that $S$ and $\psi$ are our coexisting two-particle dark matter candidates. Finally $H$ and $\Phi$ denote the $SU(2)_L$ doublet and singlet (Higgs) scalar fields, and $m_\phi$ is assumed to be smaller than 1 GeV to realize the Sommerfeld enhancement in indirect detection [13–15]. The effective scalar potential including one-loop corrections is given by

$$
V_{\text{eff}} = -\frac{1}{2} m_\phi^4 \psi^2 + \frac{\lambda_\psi}{4} \psi^4 + \lambda_1 H^\dagger H \psi^2 - \frac{1}{2} m_\phi^2 H H^\dagger H
$$

$$
+ \frac{1}{4} \lambda_1 H^\dagger H H H^\dagger H - \frac{1}{2} m_\phi^2 \Phi^2 + \frac{\lambda_2}{4} \Phi^4 + \lambda_3 \psi^2 \Phi^2 + \lambda_4 H^\dagger H H \Phi^2 + \frac{1}{64 \pi^2} \left[ m_\phi^4 \left( \ln \frac{m_\phi^2}{\mu^2} - \frac{3}{2} \right) - 3 m_\phi^4 \right]
$$

$$
+ 2 m_\phi^4 \left( \ln \frac{m_\phi^2}{\mu^2} + \frac{5}{6} \right) + 4 m_\psi \left( \ln \frac{m_\psi^2}{\mu^2} + \frac{5}{6} \right)
$$

$$
- 12 m_\phi^4 \left( \ln \frac{m_\phi^2}{\mu^2} + \frac{3}{2} \right) + m_\psi^4 \left( \ln \frac{m_\psi^2}{\mu^2} + \frac{3}{2} \right)
$$

(2)

where we have adopted $M_S$ renormalization scheme and the field-dependent masses are

$$
m^2_\psi = y_1^2 h^2 / 2,
$$

$$
m^2_\phi = (g^2 + g'^2) h^2 / 4,
$$

$$
m^2_W = g^2 h^2 / 4,
$$

$$
m^2_\phi = m^2_\phi - 2 \lambda_1 H H^\dagger H - 2 \lambda_2 \Phi^2,
$$

$$
m^2_\phi = m^2_\phi - \lambda_2 \Phi^2 - 2 \lambda_4 H^\dagger H H,
$$

$$
m^2_H = m^2_\psi - \lambda_1 H H^\dagger H - 2 \lambda_4 \Phi^2,
$$

$$
m^2_\phi = y_1^2 \Phi^2,
$$

where $\sqrt{2} H (h, 0)$. In order to guarantee the stability of the 2DM candidates, we impose the discrete symmetry $Z_2 \times Z_2^\prime$ under which all SM bosons (photon, Higgs, $W^\pm$ and $Z$) and $\Phi$ are (+, +), all SM fermions are (−, −), the singlet neutrino $S$ is (−, +) and the singlet scalar boson $\psi$ is (+, −). Now, we demand that the minimum of the scalar potential is bounded from below so as to guarantee the existence of vacuum and the minimum of the scalar potential must spontaneously break the electroweak gauge group, $(H^\dagger)$, $(\Phi) \neq 0$, but must not break $Z_2 \times Z_2^\prime$ symmetry imposed above.

Since Eq. (2) depends on the renormalization scale $\mu$, it must be RG-improved and this can be simply done by repeatedly decoupling all singlet particles and top quark at their mass scales [16]. After spontaneous symmetry breaking, the low energy scalar potential becomes

$$
V_{\text{eff}} = -\frac{1}{2} m_\psi^2 \psi^2 - \frac{1}{2} m_h^2 h^2 - \frac{1}{2} m_\phi^2 \phi^2 + 2 \phi \mu N h \phi + \frac{\lambda_1}{4} \psi^4
$$

$$
+ \frac{\lambda_1}{4} h^4 + \frac{\lambda_2}{4} \phi^4 + \frac{\lambda_3}{2} \psi^2 \phi^2 + \frac{\lambda_4}{2} h^2 \phi^2
$$

$$
+ \frac{\lambda_5}{2} \psi h^2 \phi + \frac{\lambda_6}{2} \phi h^2 \phi + \text{h.c.},
$$

(3)

where $m_\psi^2 = m^2_\phi + 2 \lambda_3 \psi^2$, $m_h^2 = \frac{1}{2} m_\phi^2 - 3 \lambda_1 v^2 - \lambda_4 v^2$, $m_\phi^2 = m^2_\phi - 3 \lambda_3 \phi^2$. Here, we have shifted the Higgs boson $H$ and the singlet Higgs scalar $\Phi$ by $H \rightarrow h + v_h$ and $\Phi \rightarrow \phi + v_\phi$, respectively, and

$$
\lambda_1 = \lambda_1 - \frac{3}{32 \pi^2} \lambda_1^2 + \frac{9}{32 \pi^2} \lambda_1^2 - \frac{3}{8 \pi^2} \lambda_2 - \frac{3}{4 \pi^2} \lambda_3 - \frac{3}{4 \pi^2} \lambda_4,
$$

$$
\lambda_2 = \lambda_2 - \frac{3}{32 \pi^2} (4 \lambda_1^2 + 4 \lambda_2^2 + \lambda_3^2 - 4 \lambda_4^2)
$$

and

$$
\lambda_4 = \lambda_4 - \frac{3}{128 \pi^2} (4 \lambda_1^2 + 8 \lambda_3^2 + 4 \lambda_2 - 4 \lambda_4).
$$

Since there exists a mixing mass term between $h$ and $\phi$, we rotate them with $\phi = \sin \theta' c \phi'$ and $h = \cos \theta' - s \phi'$, where $s$ and $c$ are $\sin \theta$ and $\cos \theta$, respectively.

For $m_\phi \lesssim 1$ GeV and $m_\phi \gg m_\phi$, the singlet neutrinos $S$ annihilate into mostly $\phi$. Other annihilation channel like $SS \rightarrow \psi \psi$ is negligible due to its very small coupling of the process. The $\phi$ can then subsequently decay into SM particles, which arises due to weak-scale WIMP annihilation as shown in [14,15].

The amount of cold dark matter in the Universe, which has been determined precisely from 5 year WMAP data [17], is given
by $\Omega_{\text{CDM}} h^2 = 0.1099 \pm 0.0062$. Assuming the coexistence of two dark matter candidates, the relic abundance observed must be composed of the contributions of both $S$ and $\psi$, $\Omega_{S} h^2 + \Omega_{\psi} h^2 = \Omega_{\text{CDM}} h^2$. The relic density of each dark matter species is approximately given by $\Omega_i h^2 \approx (0.1 \text{pb})/(\sigma v_i) \ (i = S, \psi)$, where $\langle \sigma v_i \rangle$ is the thermally averaged product of its annihilation cross section with its velocity. For our convenience, we define the parameter $\varepsilon_i$ as a ratio of $\Omega_i h^2$ to $\Omega_{\text{CDM}} h^2$.

$$\varepsilon_i = \frac{\Omega_i h^2}{\Omega_{\text{CDM}} h^2},$$

where $\varepsilon_S + \varepsilon_\psi = 1$. In fact, the parameter $\varepsilon_i$ represents the fraction of the mass density of each dark matter species in our local dark-matter halo as well as in the Universe. Since the values of $\varepsilon_i$ are unknown, we consider a few cases by choosing their values in the analysis. Each $\Omega_i h^2$ can be calculated with the help of the micrOMEGAs 2.0.7 program [18] by taking input parameters appropriately.

Except for the SM parameters, our model contains 18 new parameters: 6 scalar couplings $\lambda$, $\lambda_i$, $\lambda_{ii}(i=1\sim 4)$, 4 masses of singlet particles $M_{S,\phi,S,N}$, 3 Yukawa couplings $Y_D, S, \phi$, and 5 other parameters, $\tan \theta$, $v_{\psi,h}, \varepsilon_S, \varepsilon_\psi$. Among them, $M_N, Y_D$ and $Y_S$ are closely associated with low scale leptogenesis and light neutrino mass spectrum. There also exist 7 conditions with which parameters should be satisfied, e.g. $\varepsilon_S + \varepsilon_\psi = 1$, $c_{Vh} - 3 v_\phi = v_{EW} = 246$ GeV, etc., and that the parameters $\lambda$, $\lambda_i$, $v_{\psi,h}$ are correlated with mass parameters for $h$ and $\phi$ given in Eq. (3). Accordingly, we have 8 free parameters: the parameters $m_{S,h,\phi}$, $\lambda_3, \tan \theta$ and $\varepsilon_S$ (or $\varepsilon_\psi$) are fixed by hand and $\lambda$ (or $\lambda_3$) is determined by the conditions. (Another free parameter $\lambda_3$ is irrelevant in our analysis.) In our numerical analysis, we take $\varepsilon_\psi$ as an input parameter and then $\varepsilon_S$ is determined from the former relations and conditions. Since both $\lambda$ and $\lambda_3$ are related to the $\varepsilon_\psi$, we can take either $\lambda_3$ or $\lambda$ as an input parameter and then the other one is determined from the correlation among $\lambda$, $\lambda_3$ and $\varepsilon_\psi$.

2. Implication for the low mass DM signals

In order to interpret the low mass DM signals in terms of DM-nucleon scattering, we choose $\psi$, the lighter DM particle of 2DM, to be relevant for the experiment. Note that the heavier DM $S$ of order a few TeV is also demanded in order to explain the high energy cosmic ray anomalies later. To investigate the implication for the DM signals observed from the direct detections, we first have to estimate the DM-nucleon elastic scattering cross section predicted in our scenario. So far most experimental limits of the direct detections have been given in terms of the scattering cross section per nucleon under the assumption that there exists only one DM candidate. In the scenario of 2DM, the cross section for the WIMP-nucleon elastic scattering $\sigma_{\psi}$ is composed of $\sigma_S$ and $\sigma_\psi$ [19];

$$\frac{\sigma_S}{m_0} = \frac{\sigma_S}{m_S} + \frac{\varepsilon_\psi}{m_\psi} \sigma_\psi,$$

with $m_0$ being the WIMP mass, where we set $m_0 = m_\psi$ as the relevant DM mass for direct searches.

In our model, the non-relativistic $S$-nucleon elastic scattering cross section is given by

$$\sigma_S(\text{nucleon}) \approx \frac{1}{4\pi} \left[ \sin 2\theta Y_0 m_S^2 f \right]^2 \left[ \frac{1}{m_h^2 - m_S^2} + \frac{1}{m_\psi^2} \right],$$

where $m_h$ is a nucleon mass and $f$ is defined by the relation $f m_h = (n m_n + m_q q/3)$ whose size is determined by [20], $0.13 \lesssim f \lesssim 0.62$. The first and second terms in the parentheses correspond to the elastic scattering mediated by the Higgs field $h$ and $SU(2)_L$ singlet scalar field $\phi$, respectively. In the case of scalar $\psi$-nucleon elastic scattering, the non-relativistic elastic scattering cross section for $\psi$ is given by

$$\sigma_{\psi}(\text{nucleon}) \approx \frac{1}{4\pi} \left[ \frac{m_\psi^2 f}{m_\psi^2 + m_\psi^2} \right] \left[ \frac{m_\psi^2 f}{m_\psi^2 + m_\psi^2} \right] + \frac{\varepsilon_\psi}{m_\psi^2} \sigma_{\psi},$$

Fig. 1. (a) $\Omega_{\text{CDM}}$-nucleon elastic scattering cross section vs. DM mass. The DAMA results are presented by the grey-colored regions, which includes both channeled and quenched events as well as the allowed region consistent with the DAMA modulation signal at $3\sigma$ [21]. (In fact, the recent estimates of the channeling effect show that it is smaller than expected [22].) Also the pink-colored rectangular region corresponds to the prediction of our scenario for given input values presented in the panel and $3 \text{ GeV} \lesssim m_\psi \lesssim 11$ GeV. The red contoured region represents the DM signal from CoGeNT. (b) Allowed region of the parameter space ($\tan \theta, \varepsilon_\psi$) from the fit to the DAMA results combined with the other null experiments for $m_\psi = 3$ GeV, $\varepsilon = 0.36$ and the same input parameters as in (a). (c) Allowed region from the fit to the results for DM signal from CoGeNT for $m_\psi = 8$ GeV, $\varepsilon = 0.15$. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)
where $\lambda' = \lambda v_\phi c + 2\lambda_3 v_\phi S$ and $\lambda'' = -\lambda v_\phi c + 2\lambda_3 v_\phi c$.

In Fig. 1(a), the pink-colored rectangular area presents the predicted region of the parameter space ($\sigma_{el} - m_\phi$) in our model for several fixed input parameters given in the panel. Here, we restricted the region of $m_\phi$ to be $3 \, \text{GeV} \lesssim m_\phi \lesssim 11 \, \text{GeV}$. We see that DAMA experimental result is consistent with other null experimental results including CoGeNT 2010 (ignoring DM signal) and XENON100 data only for the narrow range $m_\phi \sim 3 \, \text{GeV}$. We also see that our predicted region for DM mass range, $7 \, \text{GeV} \lesssim m_\phi \lesssim 11 \, \text{GeV}$, is consistent with the DM signal observed from CoGeNT which corresponds to the red contour in Fig. 1(a). Fig. 1(b) represents the allowed regions of the parameter space ($\tan \theta, \lambda$) from the fit to the DAMA results combined with the other null results of direct searches particularly for $m_\phi = 3 \, \text{GeV}$ and $f = 0.36$ and the same input parameters as in Fig. 1(a). Fig. 1(c) represents the allowed parameter region from the fit to the results of DM signals from CoGeNT for $m_\phi = 8 \, \text{GeV}$ and $f = 0.15$. When we calculate numerically scattering cross sections, we vary $\lambda_3$ and $\tan \theta$ for a fixed value of $\phi S$ as well as $\lambda_\phi \lambda$ and other mass parameters of the singlet particles including Higgs boson. And then the value of $\lambda$, which lead to the right values of the scattering cross sections, can be determined accordingly. From our numerical calculation, we found that the lowest value of $\lambda$ is 0.01 which corresponds to $\lambda_3 = 0$. The allowed values of $\lambda$ increase with $\lambda_3$, but there exists the upper bound on $\lambda$ for which the scattering cross section reaches the maximally allowed value for the DM signal from CoGeNT, as can be seen in Fig. 1(c). We also notice that the excluded region $\tan \theta < 0.0022$ is not consistent with the CoGeNT DM signal because the DM-nucleon cross section size can be larger than the CoGeNT upper limit in the red contoured region. On the other hand, $\tan \theta < 0.0015$ is not consistent with electroweak symmetry breaking and relevant mass scale of the singlet scalar field $\phi$.

3. Implication for PAMELA, ATIC and Fermi-LAT

Now, let us show that the PAMELA, ATIC and Fermi-LAT data can be accounted for by regarding singlet fermion $S$ as a relevant dark matter much heavier than $\psi$, which annihilates dominantly into $\phi$, and then the $\phi$'s subsequently decay into mostly $\mu^+ \mu^-$ when $m_\phi$ is taken to be 0.25 GeV. In order to calculate the galactic cosmic ray (CR) propagation, we use GALPROP program [23] which simulates the propagation of both cosmic rays and DM annihilation products in the galaxy. The propagation equation for all CR species, can be determined accordingly. From our numerical calculation, we found that the lowest value of $\lambda$ is 0.01 which corresponds to $\lambda_3 = 0$. The allowed values of $\lambda$ increase with $\lambda_3$, but there exists the upper bound on $\lambda$ for which the scattering cross section reaches the maximally allowed value for the DM signal from CoGeNT, as can be seen in Fig. 1(c). We also notice that the excluded region $\tan \theta < 0.0022$ is not consistent with the CoGeNT DM signal because the DM-nucleon cross section size can be larger than the CoGeNT upper limit in the red contoured region. On the other hand, $\tan \theta < 0.0015$ is not consistent with electroweak symmetry breaking and relevant mass scale of the singlet scalar field $\phi$.

In Fig. 2, we present the predictions of our scenario for (a) the ratio of positron to electron plus positron fluxes and (b) the total electron plus positron fluxes, which are originated from $S \rightarrow \phi \phi$, and subsequent decays of the $\phi$'s into $\mu^+ \mu^-$ for the same input values of the parameters. As for input values, we take $\epsilon_S = 0.9$, $m_S = 2.5 \, \text{TeV}$, $m_\mu = 0.25 \, \text{GeV}$, and $(\sigma v) = 4.56 \times 10^{-26} \, \text{cm}^2 \, \text{s}^{-1}$ which satisfies the thermal relic density of $S$ for the given $\epsilon_S$. Then, the contribution of $S$ to the local DM density ($\rho_S^0$) is 0.27 GeV cm$^{-3}$ while that of $\psi$ to the local DM density ($\rho_\psi^0$) is 0.03 GeV cm$^{-3}$. In these estimates, we invoke the boost factor ($B$) reflecting Sommerfeld enhancement through which the halo annihilation rate is enhanced, and it is given by

$$B \sim \frac{\alpha m_S^2}{m_\phi},$$

where $\alpha$ lies between $10^{-3}$ and $10^{-1}$ [15]. The each curve in Fig. 2 corresponds to different boost factor, $B = 2325$ (1730, 900) for black solid (red dashed, blue solid) curve. The red dots with error bar correspond to the measurements from the PAMELA (Fig. 2(a)). The pink (grey), blue (black) circles, and the points denoted by “x” and “o” in Fig. 2(b) correspond to the measurements from ATIC, PPB-BETS, Fermi-LAT 2009 and Fermi-LAT 2010, respectively. As one can see from Fig. 2(a), the black solid and red dashed curves corresponding to $B = 2325$ and $B = 1730$ give acceptable fits to the PAMELA data for the positron fraction. For the same values of $B$, as in Fig. 2(a), the predictions of $E^3 dN/dE$ as a function of $E$ appear to give acceptable fits to the ATIC and the PPB-BETS data as well as the Fermi-LAT data for $E \lesssim 70 \, \text{GeV}$, whereas the predictions for $70 \lesssim E \lesssim 750 \, \text{GeV}$ are much deviated from the Fermi-LAT data. On the other hand, we see that the prediction of $E^3 dN/dE$ for $B = 900$ (blue curve in Fig. 2(b)) gives acceptable fit to the Fermi-LAT data, but that of positron fraction is quite small to fit
well the PAMELA data as shown in Fig. 2(a). Therefore, it looks rather difficult to perfectly accommodate the PAMELA, ATIC and Fermi-LAT data simultaneously, which may imply that there exist other astronomical sources [26–29]. In passing, please note that the introduction of one more generation of singlet vector-like neutrino, S, slightly weakens the tension between the cosmic ray data and allow for lower boost factor (even below 700). However the model loses some of its predictive power due to several new additional parameters.

In conclusion, we have shown that the extended seesaw model proposed to achieve low scale leptogenesis can resolve the anomalies in the indirect detections of annihilation products observed from PAMELA, ATIC and or Fermi-LAT and simultaneously accommodate some of recent signals of low mass DM measured at DAMA and CoGeNT. In this model, an extra vector-like singlet neutrino S and a singlet light scalar field ψ, which are coexisting two-particle dark matter candidates, are responsible for the origin of the excess positron and electron fluxes and the low mass DM signals observed from DAMA and CoGeNT. Furthermore, it has been shown that the DM signal observed from DAMA and the other null results including CoGeNT 2010 and XENON100 data from direct searches for DM can be reconciled in the case of mψ ∼ 3 GeV. We have also shown that the DM signals observed from CoGeNT can be accommodated in our model if 7 GeV ≤ mψ ≤ 11 GeV. On the other hand, in addition to SU(2)L doublet Higgs field H, the (light) singlet Higgs field Φ, which is demanded to successfully construct the coexisting two-particle dark matter scenario and whose mass is taken to be just below 1 GeV, may play an essential role in resolving the PAMELA, ATIC and or Fermi-LAT anomalies.

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