Rapid internal contraction boosts DNA friction

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Macroscopic objects are usually manipulated by force and observed with light. On the nanoscale, however, this is often done oppositely: individual macromolecules are manipulated by light and monitored with force. This procedure, which is the basis of single-molecule force spectroscopy, has led to much of our quantitative understanding of how DNA works, and is now routinely applied to explore molecular structure and interactions, DNA–protein reactions and protein folding. Here we develop the technique further by introducing a dynamic force spectroscopy set-up for a non-invasive inspection of the tension dynamics in a taut strand of DNA. The internal contraction after a sudden release of the molecule is shown to give rise to a drastically enhanced viscous friction, as revealed by the slow relaxation of an attached colloidal tracer. Our systematic theory explains the data quantitatively and provides a powerful tool for the rational design of new dynamic force spectroscopy assays.
Life literally is a tour de force. And much of the mechanical load exerted on tissues and cells is supported and transmitted by biomolecules that can be classified as semiflexible polymers. Already at physiological force levels, semiflexible biopolymers exhibit non-Hookean stretching behaviour, which sets them apart from other nanoscopic force-bearing elements, such as AFM tips or optical tweezers. Their characteristic force-extension relation under static tension has become the workhorse of a whole industry of single-molecule force spectroscopy methods, employing the polymers as gauges, linkers, molecular handles, or as a virtual magnifying glass to explore the molecular world at ever higher spatial and temporal resolution. However, the mechanical nonlinearity substantially complicates the dynamic response to external perturbations. If you excite an overdamped linear spring, the excitation always decays exponentially and the spring attains its equilibrium state within a characteristic relaxation time. In contrast, polymeric response curves exhibit no well-defined relaxation timescale and their relaxation generally passes through a multitude of power-law regimes, subtly dependent on the details of the applied force protocol. Mathematically speaking, the response is not governed by the relaxation of a single eigenmode (a single equivalent spring-dashpot element) but by a whole hierarchy of bending fluctuations of different wavelengths and their corresponding relaxation times. As demonstrated below, such complex mode superpositions can give rise to somewhat unintuitive and sometimes surprising dynamical effects.

A major source of complication in dynamics, compared with the better understood stationary situation, lies within the intricate coupling of the conformational dynamics of a stretched polymer to the solvent hydrodynamics. In its equilibrium conformation, as a coil of radius R, the polymer has a friction coefficient comparable to that of a colloidal bead of about the same radius. In contrast, if the polymer has a straight, rod-like conformation, it feels a friction force proportional to its contour length L. For a λ-phage DNA, L/R ≈ 20, which underscores how important the effect can become. The result of a more quantitative theoretical analysis for our specific experimental set-up is illustrated in Fig. 1. It shows that the friction coefficient of a DNA-coated colloidal bead (Fig. 1b) would increase by ≈30% if one of the molecules of its coating layer could be permanently transformed into a straight conformation of length L = 16 μm (Fig. 1c). The purpose of the following discussion is to demonstrate both experimentally and theoretically that a much more dramatic, roughly fourfold, friction enhancement can be achieved if the initially stretched polymer is additionally allowed to recoil dynamically (Fig. 1d). At the origin of this observation lies an excess viscous friction that is self-generated by the internal conformational dynamics. To our knowledge, this is the first time this effect has been isolated experimentally, as earlier measurements were either limited to small fluctuations around the stretched state or suppressed the generation of excess friction through the attachment of large beads to the recoiling polymer end. In the following, we elucidate the underlying physics of the enhanced friction and describe our dedicated set-up to measure it. We then present our measurement results and interpret them in the light of a systematic theory of nonequilibrium polymer dynamics.

In semiflexible polymer dynamics, hydrodynamic friction arises both due to the polymer's transverse bending undulations and due to the contour's longitudinal motion with respect to the background solvent. For a weakly bending (stretched) polymer, longitudinal fluctuations are geometrically strongly suppressed relative to transverse fluctuations. The linear dynamics of transverse bending modes therefore not only dominates moderate transverse contour excursions, but also slow longitudinal fluctuations—for example, the dynamic longitudinal mean-square displacement of a monomer at late times t → ∞. But theory predicts that the short-time longitudinal dynamics, which becomes accessible by our dedicated experimental set-up, is strongly nonlinear and dominated by longitudinal drag. An example is provided by the initial stretching dynamics in response to a suddenly applied pulling force. But such a fast stretching is limited by longitudinal drag for a rigid-rod scenario (DNA-coated bead with one DNA molecule straightened out): γ ≈ 2πηL/\log(L/d) + γ_{bead} ≈ 39 pN ms μm⁻¹, τ ≈ 1.7 ms. (d) Stretch relaxation (DNA-coated bead with one initially stretched polymer that is freely retracting); γ ≈ 39 pN ms μm⁻¹, τ ≈ 5 ms (determined via f_{trap}(t) = f_{trap}(t)/e(t)). (e) The recoiling of a polymer from a stretched initial state is entirely driven by entropic forces or 'thermal noise' (red arrows). The additional contour length required to create the wiggles in the contour has to be pulled in from the polymer ends. The concomitant solvent flow (blue) exerts counteractive longitudinal drag forces that limit the maximum retraction speed and ultimately lead to the excess friction seen in scenario (d).
allows for the best experimental control over the comparatively scarcely investigated longitudinal motion limited by longitudinal drag, as it gives rise to an extended phase of quasistatic homogeneous tension relaxation limited by longitudinal drag\(^1,21\). A first glimpse of this type of motion—a power-law growth of the retraction or ‘length deficit’ with \(t^{1/3}\)—has already previously been caught\(^22\). But with our set-up, which enables us to directly monitor the decay of the backbone tension, we can for the first time inspect the underlying physical mechanism. As we demonstrate next, the recoiling, although it is solely driven by entropy, \(\nu(z)\), by the thermal forces resulting from the random bombardment of the polymer backbone with solvent molecules, involves such high local retraction velocities as to create the massive excess friction alluded to above (Fig. 1d).

**Results**

**Measurements.** Our dedicated set-up to reversibly and non-invasively stretch and release a strand of \(\lambda\)-phage DNA is sketched in Fig. 2. A comprehensive description of the experimental procedures can be found in the Methods. Briefly, a calibrated optical trap is embedded within a microfluidic cell serves to hold or relocate a single DNA-coated colloidal bead\(^23\). One of the grafted DNA molecules on the bead is pulled straight by applying an electrophoretic force with a nearby nanocapillary. After the electric field has abruptly been switched off at time \(t = 0\), a high-speed video camera tracks the bead displacement in the trap with submillisecond time resolution\(^24\). From the bead position, we infer the force \(f_{\text{trap}}(t)\) exerted by the trap onto the bead. Up to the relatively small bead friction (that is fully included in our theory), \(f_{\text{trap}}(t)\) is the value \(f_\nu(L, t)\) of the backbone tension profile \(f_\nu(s, t)\) at the grafted end \(s = L\). Typical time traces of the force \(f_{\text{trap}}(t)\) exerted by the trap onto the bead are depicted in Fig. 3.

**Theory.** Our systematic theory is in good quantitative agreement with the data (Fig. 3, lines). Since the optically trapped bead moves by no more than \(\approx 200\,\text{nm}\) during the whole experiment, the bead-attached end can be thought fixed, thus closely approximating the “release” scenario described in refs. 16, 25. Yet, taking also the slight bead motion into account, the theory allows us to reconstruct from \(f_{\text{trap}}(t)\) the complete spatio-temporal evolution of the tension profile \(f_\nu(s, t)\), including the tension \(f_\nu(L, t)\) at the grafted polymer end, (Fig. 4) and the local retraction...
which the growth and decay of the retraction velocity along the polymer can be inferred\(^{10,11,16,25}\), as depicted in Fig. 5.

**Discussion**

The exact mechanism giving rise to the observed slow relaxation behaviour is intuitively best understood as a dynamical equivalent of the static worm-like chain force-extension relation. On a scaling level, we balance the effective retraction force \(f_\xi - L k_T z\) (ref. 2) that retards the retraction of the end-to-end distance \(z\) of the molecule by the entropic driving force \(f_\xi = k_B T / (4 f_p (L - z)^2 /)\) that follows from the well-known Marko–Siggia expression for an almost straight conformation. Time integration then yields the aforementioned \(f_\xi^{1/3}\)—growth of the deficit \(L - z\) and the asymptotic power-law decay \(f_\xi(t) \sim \left( z^2 / k_T L / f_p \right)^{1/2} t^{2/3}\) of the tension\(^{25,26}\). By allowing both the tension and the contraction to vary along the backbone, the argument can be made rigorous to yield the correct partial differential equation governing the evolution of \(f_\xi(s, t)\) (refs 21,25). Further, including the optical trap and the bead friction, we obtain a slightly more complicated set of equations that can be solved numerically. To ease data analysis in future applications, we have extracted an approximate semi-empirical formula from this first-principles theory (the relevant equations of motion and the definitions of \(\chi\) and \(\alpha\) are given in the Methods section),

\[
\begin{align*}
 f_p(t) & \approx \frac{\beta f_\xi^0}{1 + 9(t/t_0)^2} + (1 - \beta) f_\xi^0 e^{-t/(k_B t)} \\
 t_0 & = \frac{\xi_1}{4} \left( k_B T k_T / f_p \right)^1/2 \\
 \beta & \approx \frac{\chi / z}{\chi / z + 0.6(z / \chi)^{0.56}} + \frac{1}{1 + z} \\
 \tilde{v} & \approx \frac{1 + \beta}{\chi / z} + \frac{1}{1 + z} + \frac{0.16}{z / \chi} 
\end{align*}
\]

This formula was validated against the exact numerical solutions which it matches closely over a wide range of polymer lengths \(L\), bead sizes \(\gamma_{\text{head}}\), trap stiffnesses \(k_T\), persistence lengths \(L_p\) and initial values of the tension \(f_\xi^0\). It interpolates between the purely exponential retraction obtained for negligible excess friction (as, for example, in the limit of a very short/stiff polymer) and the algebraic retraction \(f_\xi(t) \sim t^{-2/3}\) that would result for a rigidly fixed polymer end\(^{25}\) (corresponding to the limit of a very stiff trap). Thereby, it provides a precise quantification of the friction enhancement anticipated in Fig. 1 and a very convenient starting point for further quantitative applications of our dynamic force spectroscopy setup. Both dimensionless parameters \(\chi\) and \(\alpha\) are unambiguously defined in terms of experimental conditions (cf. Methods), thus leaving no free parameter save for the effective friction coefficient \(\xi_1\). The friction coefficient per length \(\xi_0\) can be estimated asymptotically \(\xi_0 / L = 2 \pi \eta / \log(L/a) \sim 0.65 \text{ mPa s} \) (for a stretched molecule of diameter \(a = 2 \text{ nm}\) (ref. 27)), and we have used this value for all of our theory curves. This renders the comparison between data and analytical theory in Fig. 3 parameter free.

In summary, combining experiment and theory, we have developed a quantitative dynamic force spectroscopy set-up. We demonstrated that the internal dynamics of a recoiling taut DNA molecule generates a large excess friction that drastically delays the relaxation of an attached tracer bead, as quantitatively predicted by a simple but precise analytical approximation to the systematic theory. The theory also enabled us to reconstruct the complete spatio-temporal evolution of both the backbone tension and the rebound velocity along the polymer, from records of the tracer relaxation. In the future, it might be interesting to use this

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**Figure 4** | Evolution of the backbone tension along the polymer.
Numerical solutions of equation 4 for a set of three different initial conditions (corresponding to boundary layer widths of about 1, 2 and 4 \(\mu\)m coloured cyan, blue (dashed) and black (dotted), respectively) at various times \(t\) (in ms). The tension profiles quickly converge from their different initial states and settle onto a unique scaling function \(f_p(s)\) with a slowly decaying amplitude, rendering the tension at the bead-laden end independent of the precise initial condition.

**Figure 5** | Retraction velocities along the polymer backbone at different times.
Theoretically reconstructed profiles \(v(s, t)\) of the local longitudinal contour velocity, based on the dynamic decay of the force \(f_{\text{trap}}(t)\) on the bead in Fig. 3. The computed velocity profile quickly settles down to an almost linear profile after the propagation layer has swept through the polymer (at \(t \approx 1\) ms). Yet, substantial retraction velocities persist near the free end, and the resulting enhanced friction (Fig. 1c) considerably retards the decay of the backbone tension at the grafted end. Inset: qualitative sketch (not to scale) of the recoil process and the accompanying change in backbone tension (from top to bottom).
quantitative set-up to reveal how the characteristic dynamics is modified by the presence of DNA–protein interactions, nicks or supercoils.28

Methods

Experimental. The microfluidic cell consists of polydimethylsiloxane (PDMS) and encompasses two 100 μl reservoirs connected by a nanocapillary, whose orifice size defines the active sensing volume for single-molecule studies. We use quartz glass capillaries (Hilgenberg, Germany) with an outer diameter of 0.5 mm and a wall thickness of 0.064 mm (20 mM) and 0.1 mm (50–750 mM), respectively. At the bottom the microfluidic cell is sealed with a glass cover slide of 100 μm thickness. The electric potential is applied with the help of two Ag/AgCl electrodes, one of them residing within the nanocapillary, the other located in front of the capillary tip. We use a potassium chloride (KCl) solution at concentrations of 20, 50, 100, 500 and 750 mM in 10 mM Tris buffer at pH 8.

Our custom-built optical tweezers set-up is assembled on an optical table and uses a 5-W ytterbium fibre laser (YLM-5-LP, IPG Laser, Germany) at a wavelength of λ = 1.064 μm (ref 24). Whereas the optical trap itself is static, the microfluidic cell is mounted onto a xyz piezo nanopositioning system (P-517.3 and E-710.3, Physik Instrumente, Germany). With a range of 100 μm in x- and 20 μm in z-direction it allows for nm-precise manoeuvring of entrapped beads in relation to the surrounding cell. Illumination of the region of interest inside the sample cell is either done using a standard white light source (DC-950 Fiber-Lite, Edmund Optics, USA), or an optical fibre (100 W mercury arc lamp, LOT-Oriel, UK and 600 μm multimode silica fibre, NA0.39, Thorlabs, UK). In combination with a high-speed CMOS camera (MC1362, Mikrotron, Germany) this modular approach allows for real-time tracking of optically trapped colloids at up to 10,000 fps (frames per second) with 2 nm spatial resolution.24

A Faraday cage encloses the sample cell. The externally applied electric potential is held constant by a commercial amplifier (Axopatch 200B, Axon Instruments, USA) in voltage-clamp mode. Its headstage is mounted inside the Faraday cage and connected to the Ag/AgCl electrodes of the microfluidic cell. Concurrent measurements of the ionic current through the nanocapillary allow the simultaneous trapping of multiple DNA strands to be ruled out. Before each experiment we determine the nanocapillary tip diameter via its current–voltage characteristics.23 Afterwards, λ-DNA-coated colloids are flushed into the microfluidic cell and the optical trap is calibrated by analysing its power spectral density.24,32. The proper grafting of λ-DNA to our polystyrene colloids is verified by repeating the power spectrum calibration for a number of particles.

Tension dynamics. Apart from a short-lived initial regime (t ~ ns) that we cannot measure here, the relaxation behaviour of a freely contracting semiflexible polymer is governed by a generalization of the Marko–Siggia force–extension relation2 to space- and time-dependent backbone tension $f_P(s, t)$.

\[
\begin{align*}
\partial_t^2 f_P(s, t) &= \frac{1}{4} \sqrt{\frac{k_B T}{f_P(s, t)}} \partial_s f_P(s, t),
\end{align*}
\]
For a freely relaxing chain, the above equation has been derived and solved systematically in the literature.\(^6\)\(^\text{–}\)\(^8\) Within this theoretical framework, our stretching apparatus constitutes an initial condition \(f_p(x, t = 0)\), where the force gradient \(\partial f_p(x, t = 0)\) is linearly proportional to the electric field and thus to the resistivity determined by the nanocapillary cross section,

\[
\partial f_p(s, t = 0) \propto E(s) \propto \frac{1}{\eta_{\text{cap}}(s)^{\frac{3}{2}}}
\]  

(3)

The fairly complex distribution of forces acting within the nanocapillary provides a good experimental approximation to a point force applied to the polymer end, as details of the initially applied tension profile \(f_p(s)\) quickly diffuse away during tension propagation, see Fig. 4. Although in principle equation 2 breaks down close to the force-free end, and hydrodynamic interactions with the pore entrance or hydrodynamic boundary effects might locally increase the longitudinal friction coefficient \(\zeta_L\); these finer points turn out to be negligible (Supplementary Notes 1 and 2, Supplementary Figs S3 and S4).

We now extend equation 2 to dynamic force spectroscopy assays by coupling one polymer end to a linear trap of stiffness \(k_{\text{trap}}\) and drag coefficient \(b_{\text{trap}}\). Given a certain time-dependent bead velocity \(v_{\text{head}}(t)\) that must agree with the local polymer velocity \(\partial f_p(s = L, t)/\partial s\) at the bead-laden end, the dynamic tension profile is then determined as follows,

\[
f_p(s, 0) = f_p(s) \quad f_p(0, t = 0) = 0
\]

\[
\partial f_p(L, t)/\partial s = v_{\text{head}}
\]

\[
\partial^2 f_p(s, t) = \frac{1}{4} \frac{\zeta L}{k_B T} \partial f_p(s, t) \frac{g_t}{f_p(s, t)^{\frac{3}{2}}}
\]

(4)

The bead velocity \(v_{\text{head}}\) is not an external parameter but depends on the polymeric force \(f_p\) exerted on the bead. As bead and trap together constitute a linear overdamped subsystem with characteristic relaxation time \(\eta_{\text{trap}}/b_{\text{trap}}\), we can infer the bead velocity from the time-dependent polymer backbone tension evaluated at its bead-attached end \(f_p(L, t)\),

\[
v_{\text{head}}(t) = -\frac{1}{\eta_{\text{trap}}/b_{\text{trap}}} \int f_p(L, t') e^{-\frac{b_{\text{trap}}}{k_B T} t - \tau} dt'.
\]

(5)

This closes equation 4, which we solve self-consistently by iterating in \(v_{\text{head}}(t)\) (Supplementary Methods).

To obtain a uniformly valid practical approximation to the resulting force relaxation \(f_p(t) := f_p(L, t)\), we first reexpress the above system of equations in dimensionless form by measuring \(f_p\) in units of \(f^\text{th}\), distance \(s\) in units of polymer length \(L\) and time \(t\) in units of \(t_b \equiv \frac{\zeta L^2}{4f_p^2} \sqrt{k_B T/f_p}\).

The equations of motion then read

\[
\partial f_p(1, t) = -\alpha \int_0^1 f_p(s, t') e^{-\alpha(t-t')} ds' 
\]

\[
\partial^2 f_p(s, t) = \frac{\zeta}{f_p(s, t)} \frac{g_t}{f_p(s, t)^{\frac{3}{2}}}
\]

where

\[
\alpha \equiv \frac{\zeta L}{b_{\text{trap}}}
\]

(6)

denotes the ratio between ‘static’ polymer drag and bead drag and

\[
\chi \equiv \frac{\chi_{\text{trap}}}{\chi_{\text{head}}^2} \frac{L^2}{4f_p^2} \frac{\sqrt{k_B T}}{b_{\text{trap}}}
\]

(7)

does not affect the form of the equations. However, this allows us to solve our data closely approaches the asymptotic solutions \(\beta(x = \infty) \approx 1\) (since \(\chi \to \infty\) can be realized in measuring \(k_{\text{trap}}\) to infinity, leading back to the pure release scenario), \(\beta(x = 0) = 1\) (directly follows from the dimensionless equations of motion) and \(\beta(2x = 0) = 1/(1 + x)\). The latter asymptote corresponds to a virtually infinite persistence length \(\eta_{\text{cap}}\gg L\) such that the polymer retracts as a solid object of drag coefficient \(\zeta L\) and producing purely exponential relaxation on a timescale \(t\) as discussed above. In contrast, the timescale of internal contraction remains unchanged in units of \(t_b\) such that tension propagation amounts to a sudden drop from the initially applied force \(f_p(t = 0)/f_p^\text{th} = 1\) to some smaller force \(f_p(0)/f_p^\text{th} = \beta\), which we identify with the force needed to drag along the polymer at an instantaneous speed \(x_0^\text{th}\). As

\[
x(t) = x_0 e^{-\chi t/\beta} = \frac{f_p^\text{th}}{k_{\text{trap}}} e^{-\chi t/\beta}.
\]

(8)

For large \(x\), we do not have an analytic expression for finite \(\chi\). However, as for any finite \(\chi\), \(\beta\) must approach zero as \(x \to \infty\), we know that all nontrivial values of \(\beta\) are realized for very large values of \(x\) as \(\chi\) diverges. The boundary condition at the bead-attached end thus simplifies to

\[
\partial f_p(1, t) \sim -\frac{2}{\chi} f_p(1, t).
\]

We thus only need to make sure that our numerical data approaches a self-similar shape \(f_p(y, x) = f_p(x/x')\) in the limit of large \(x\). Taking this into account, we arrive at the approximate expressions for \(\beta\) and \(x\) shown in equation 1. Figure 6 illustrates the quality of our empirical interpolation formula for nonasymptotic values of \(\chi\) and \(x\).

References

1. Kasza, K. et al. The cell as a material. Curr. Opin. Cell Biol. 19, 101–107 (2007).
2. M. J. F. & Siggea, E. D. Stretching DNA. Macromolecules 28, 8759–8770 (1995).
3. Bustamante, C., Bryant, Z. & Smith, S. Ten years of tension: single-molecule DNA mechanics. Nature 421, 423–427 (2003).
4. Gebhardt, J., Bornschlegl, T. & Rief, M. Full distance-resolved folding energy landscape of one single protein molecule. Proc. Natl Acad. Sci. 107, 2013–2018 (2010).
5. Aubin-Tam, M. et al. Adhesion through single peptide aptamers. J. Phys. Chem. A 115, 3657–3664 (2011).
6. Koch, S. & Wang, M. Dynamic force spectroscopy of protein-DNA Interactions by unzipping DNA. Phys. Rev. Lett. 91, 028103 (2003).
7. Li, H. et al. Reverse engineering of the giant muscle protein titin. Nature 418, 998–1002 (2002).
8. Grandbois, M., Beyer, M., Rief, M., Clausen, H. & Gaub, H. How strong is a covalent bond? Science 283, 1727–1730 (1999).
9. Müller, D. & Anderson, K. Biomolecular imaging using atomic force microscopy. Trends Biotechnol. 20, 545–549 (2002).
10. Hallatschek, O., Frey, E. & Kroy, K. Propagation and relaxation of tension in stiff polymers. Phys. Rev. Lett. 94, 077804 (2005).
11. Obermayer, B., Möbius, W., Hallatschek, O., Frey, E. & Kroy, K. Freely relaxing polymers remember how they were straightened. Phys. Rev. E 79, 021804 (2009).
12. Meiners, J. & Quake, S. Femtonewton force spectroscopy of single extended DNA molecules. Phys. Rev. Lett. 84, 5014–5017 (2000).
13. Goshen, E., Zhao, W., Carmon, G., Rosen, S. & Feinberg, M. Relaxation dynamics of a single DNA molecule. Phys. Rev. E 71, 061920 (2005).
14. Bohbot-Raviv, Y., Zhao, W., Feingold, M., Wiggins, C. & Granek, R. Relaxation Dynamics of semiflexible polymers. Phys. Rev. Lett. 92, 098101 (2004).
15. Crut, A., Koster, D. A., Seidel, R., Wiggins, C. H. & Dekker, N. H. Fast dynamics of supercoiled DNA revealed by single-molecule experiments Proc. Natl Acad. Sci. 104, 11957–11962 (2007).
16. Hallatschek, O., Frey, E. & Kroy, K. Tension dynamics in semiflexible polymers. I. Coarse-grained equations of motion. Phys. Rev. E 75, 031905 (2007).
17. Seifert, U., Wintz, W. & Nelson, P. Straightening of thermal fluctuations in semiflexible polymers by applied tension. Phys. Rev. Lett. 77, 5389–5392 (1996).
18. Ajdari, A., Jülicher, F. & Maggs, A. Pulling on a filament. J. Phys. I France 7, 823–826 (1997).
19. Everaers, R., Jülicher, F., Ajdari, A. & Maggs, A. Dynamic fluctuations of semiflexible filaments. Phys. Rev. Lett. 82, 3717–3720 (1999).
20. Turner, S., Cabodi, M. & Craighead, H. Confinement-induced entropic recoil of single DNA molecules in a nanofluidic structure. Phys. Rev. Lett. 88, 128103 (2002).
21. Brochard-Wyart, F., Buguin, A. & de Gennes, P. Dynamics of taut DNA chains. Europhys. Lett. 47, 171–174 (1999).
22. Schroeder, C., Shaqfeh, E. & Chu, S. Effect of hydrodynamic interactions on DNA dynamics in extensional flow: simulation and single molecule experiment. Macromolecules 37, 9242–9256 (2004).
23. Steinbock, L., Otto, O., Chimerel, C., Gornall, J. & Keyser, U. Detecting DNA folding with nanocapillaries. Nano Lett. 10, 2493–2497 (2010).
24. Otto, O. et al. Real-time particle tracking at 10,000 fps using optical fiber illumination. Opt. Express 18, 22722–22733 (2010).
25. Hallatschek, O., Frey, E. & Kroy, K. Tension dynamics in semiflexible polymers. II. Scaling solutions and applications. Phys. Rev. E 75, 031906 (2007).
26. Shaqfeh, E., McKinley, G., Woo, N., Nguyen, D. & Sridhar, T. On the polymer entropic force singularity and its relation to extensional stress relaxation and filament recoil. J. Rheol. 48, 209–221 (2004).
27. Batchelor, G. K. Slender-body theory for particles of arbitrary cross-section in Stokes flow. J. Fluid Mech. 44, 419–440 (1970).
28. van Loenhout, M. T. J. et al. Dynamics of DNA supercoils. Science 338, 94–97 (2012).
29. Otto, O., Steinbock, L., Wong, D., Gornall, J. & Keyser, U. Direct force and ionic-current measurements on DNA in a nanocapillary. Rev. Sci.Instrum. 82, 086102 (2011).
30. Steinbock, L. et al. Probing DNA with micro- and nanocapillaries and optical tweezers. J. Phys. Condens. Matter 22, 454113 (2010).
31. Holmberg, A. et al. The biotin-streptavidin interaction can be reversibly broken using water at elevated temperatures. Electrophoresis 26, 501–510 (2005).
32. Berg-Sörensen, K. & Flyvbjerg, H. Power spectrum analysis for optical tweezers. Rev. Sci. Inst. 75, 594–612 (2004).
33. Manning, G. S. The persistence length of DNA is reached from the persistence length of its null isomer through an internal electrostatic stretching force. Biophys. J. 91, 3607–3616 (2006).
34. Smeets, R. M. M., Keyser, U. F., Krapf, D., Wu, M. Y., Nynke, H. & Dekker, C. Salt dependence of ion transport and DNA translocation through solid-state nanopores. Nano Lett. 6, 89–95 (2006).

Acknowledgements
We thank Ralf Seidel for discussions, Benedikt Obermayer for providing us with a numerical solver for the full (non-quasistatic) theory of tension propagation and Lorenz Steinbock for the preparation of our glass nanocapillaries. S.S. acknowledges financial support from the Deutsche Forschungsgemeinschaft (DFG) through FOR 877, the European Social Fund and the Leipzig School of Natural Sciences—Building with Molecules and Nano-objects (BuildMoNa). O.O. is supported by a Ph.D. fellowship from the Boehringer Ingelheim Fonds. U.F.K. was supported through an Emmy Noether grant from the Deutsche Forschungsgemeinschaft and an ERC starting grant. N.L. acknowledges financial support from the George and Lillian Schiff Foundation and Trinity College, Cambridge.

Author contributions
U.K., O.O. and N.L. devised and carried out the measurements; O.O. and S.S. analysed the experimental data; S.S. and K.K. provided the theoretical description; K.K., U.K. and S.S. wrote the paper. O.O. and S.S. contributed equally to the study. ALL authors discussed the results and commented on the manuscript.

Additional information
Supplementary Information accompanies this paper at http://www.nature.com/naturecommunications

Competing financial interests: The authors declare no competing financial interests.

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How to cite this article: Otto, O. et al. Rapid internal contraction boosts DNA friction. Nat. Commun. 4:1780 doi: 10.1038/ncomms2790 (2013).

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