Bulk electronic transport impacts on electron transfer at conducting polymer electrode–electrolyte interfaces

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Electrochemistry is an old but still flourishing field of research due to the importance of the efficiency and kinetics of electrochemical reactions in industrial processes and (bio-)electrochemical devices. The heterogeneous electron transfer from an electrode to a reactant in the solution has been well studied for metal, semiconductor, metal oxide, and carbon electrodes. For those electrode materials, there is little correlation between the electronic transport within the electrode material and the electron transfer occurring at the interface between the electrode and the solution. Here, we investigate the heterogeneous electron transfer between a conducting polymer electrode and a redox couple in an electrolyte. As a benchmark system, we use poly(3,4-ethylenedioxythiophene) (PEDOT) and the Ferro/ferricyanide redox couple in an aqueous electrolyte. We discovered a strong correlation between the electronic transport within the PEDOT electrode and the rate of electron transfer to the organometallic molecules in solution. We attribute this to a percolation-based charge transport within the polymer electrode directly involved in the electron transfer. We show the impact of this finding by optimizing an electrochemical thermogalvanic cell that transforms a heat flux into electrical power. The power generated by the cell increased by four orders of magnitude on changing the morphology and conductivity of the polymer electrode. As all conducting polymers are recognized to have percolation transport, we believe that this is a general phenomenon for this family of conductors.

Heterogeneous electron transfer (HET) takes place at the interface between a solid electrode and a reactant in a liquid electrolyte. HET is of uttermost importance in many technologies today, such as corrosion protection, fuel cells, batteries, medical sensors, organic electrochemical synthesis, food analysis, and electronics. The rate of the HET depends on the type of reactants, electrodes, solvents, and electrolyte. Interestingly, HET has been studied on various types of planar and nonporous electrodes: liquid and solid metals, semimetals such as graphite, and inorganic semiconductors. HET can be classified in two routes: when the reactant can (cannot) chemisorb on the electrode; in that case, the HET is called “inner sphere” and fast (“outer sphere” and slow), indicating strong (weak) electronic coupling between the electrode and the reactant (1). In first approximation (2), the faradic current $i_f$ generated by the HET is proportional to the kinetic constant of the reaction $k_o$, the density of surface charge carriers of relevant energy $N_s(E)$, the concentration of reactant at the surface $C_{aq}$, and an energy distribution function related to the electronic level of the reactant-involved $E^{R}(E)$:

$$i_f = k_o N_s(E) C_{aq} W(E).$$

[1]

An electron transfer at semiconductor or semimetal electrodes is often slower than that on metals, not necessarily because of the absence of chemisorption but rather, because of a low density of electronic states $N(E)$ and a corresponding low density of charge carriers (3). The faradic current from the HET can also be enhanced with porous electrodes of the large surface area (4) with the extreme case of carbon nanotubes. Metallic carbon nanotubes have a surface area of 1,100 m²/g (5). For a wide range of potential, the nanotubes are stable and metallic, thus constituting a good electrode. Finally, there exist conducting polymer electrodes that resemble none of the existing classes of materials. However, they have unique features, making them attractive for electrocatalysis (6), fuel cells (7), batteries (8), and sensors (9).

Conducting polymers have a lower density of state typically at the Fermi level than metals, which is also reflected by their modest conduction ($1–4 \times 10^{-3}$ S/cm), approaching that of bad metals, like Hg (10⁻⁹ S/cm), but much less conducting than Cu (6 × 10⁻⁸ S/cm). Compared with the metal electrode, we can distinguish four major unique features. First, conducting polymers are stable in air and do not form an insulating oxide layer like conventional nonnoble metals. Second, when swelled in a solvent, they possess a molecular porosity, enabling some small molecular reactants to penetrate the bulk of the electrode: thus, potentially, an intrinsic very high surface area for small reactants. Third, the presence of noncovalent bonds between the polymer chains makes them favorable for ion transport. Hence, conducting polymers are classified among the mixed electron–ion conductors (10). Fourth, they are electrochemically active solids (11, 12). Hence, the electron transfer with the reactant in the conducting polymer | electron transfer | thermogalvanic cell
solution must occur at an electrochemical potential corresponding to the conducting state of the polymer electrode; otherwise, the polymer would be insulating and cannot play the role of an electrode (i.e., no efficient electron transfer can occur) (8). A recent study shows the irreversible and strong potential dependence of the rate of electron transfer on the density of state of the polymer electrode in that specific case where the energy distribution function $W^R(E)$ of the reactant is located at the band edge of the polymer electrode [that is, where the $N(E)$ drops to zero in the band gap (13)]. Finally, the chemical structure of the polymer electrode creates specific catalytic sites boosting electrocatalytic reactions, such as, for example, CO$_2$ reduction (14) and O$_2$ reduction (15). This is an exciting avenue, because the chemical design of the polymer can be optimized for a specific electrochemical reaction.

Although semimetallc (16) and metallic behaviors (17, 18) have been found in specific samples in conducting polymers with conductivity decreasing on heating, paracrystalline or amorphous doped polymers usually display a temperature-activated transport in a large temperature range (19). Conducting polymers appear as inhomogeneous solid in which the transport is limited by the connectivity of conducting ordered domains because of the presence of poorly conducting disordered domains acting as an energy barrier; the latter causes the temperature dependence of conductivity. In these materials, percolation transport is proposed as a predominant mechanism (20, 21). Poly(3,4-ethylenedioxythiophene) (PEDOT) complexed with poly(styrenesulfonic acid) is used as a benchmark system, because it is the most widely used among all intrinsically conducting polymers due to its water processability, its high electrical conductivity (up to about 1,000 S/cm), and good air stability (22). The chemical structures are presented in Fig. 1. The subvalent-to-thiophene ratio is about three, indicating a clear excess in poly(styrenesulfonate) (PSS) not involved in the neutralization of the positively charged PEDOT chains (23). The PEDOT-PSS nanoparticles are 20–30 nm in diameter in solution (24).

![Fig. 1. Schematic diagram of charge transport (pink arrows) within the polymer electrode and the electron transfer (green arrows) at the polymer electrolyte interface for (A) poor-conducting, (B) intermediate-conducting, and (C) high-conducting PEDOT-PSS electrodes. The blue "snake" lines (on left-hand side) mimic aggregates of PEDOT chains that display short-range order through $\pi-\pi$ stacks.](image)

Results and Discussion

Importantly, we do not use a metallic electrode underneath the polymer electrode to avoid any parasitic side reaction with the metal and to truly reveal the interplay between HET and electronic transport within the polymer. Our choice of PEDOT-PSS and Ferro/ferricyanide redox couple in aqueous electrolyte (Fig. 1). We investigate the interplay between the electrical conductivity of the polymer, its density of state $N(E)$, and the rate of electron transfer $\dot{E}$ to the reactant and show that an additional phenomenon is involved in the description of HET at conducting polymer electrodes. We use this knowledge to optimize the electrical power generated by a thermogalvanic cell (TGC), which is an electrochemical device that converts a heat flow into electricity.
various amounts of DMSO up to 6 wt %. The electrical conductivity measured with a four-point probe technique is reported as a function of DMSO content. A sigmoid evolution is observed starting with low conductivity (0.4 S/cm) for pristine PEDOT-PSS (0 wt % DMSO) to 805 S/cm for 5% DMSO (Fig. 2A). Our maximum conductivity (805 S/cm) is slightly lower than in previous reports (900 S/cm) (32). The high conductivity obtained for the best PEDOT-PSS morphology is explained by a high-oxidation/doping level [on average, one positive charge for three monomer units (23)] in combination with high mobility of about 0.3–2.7 cm² V⁻¹ s⁻¹ (10, 19). The surprising high mobility observed while PEDOT-PSS is rather amorphous indicates that the charge transport is governed by short-range order (33, 34) into small nanocrystalline domains, as also evidenced by the presence of π–π stacking (35) in the overall amorphous solid. The sigmoid behavior is in agreement with previous observations of a phase separation phenomenon within the PEDOT-PSS blend associated with a separation of the excess PSS (26). The amount of PEDOT metallic nanocrystals is not dramatically changed with the DMSO content as indicated by two observations. First, the absorption spectrum is practically unchanged with a large absorption background in the infrared specific to Drude-like absorption features typical of metallic domains (36) (SI Appendix, Fig. S1). Second, the surface density of state N(E) at the Fermi level measured by UV-photoclectron spectroscopy (UPS), which indicates a vanishingly small band gap, does not change despite the variation in DMSO content (Fig. 2B). While UPS is performed in vacuum, we also confirm that the electrochemical potential of the PEDOT-PSS electrode does not change with the DMSO content when it is swelled with an aqueous electrolyte (SI Appendix, Fig. S2). Hence, the three-orders of magnitude increase in conductivity is attributed to the formation of percolation paths realized by a reorganization of the morphology at the nanoscale to promote better connectivity between the metallic-like nanocrystals (10, 20, 21). The model proposed in the literature is in agreement with what is found here despite the addition of GOPS. Fig. 1 shows the formation of percolation paths within the PEDOT electrode. The better conductivity is indicated by a lowering of the activation energy for the transport extracted from a simple Arrhenius law of the conductivity vs. temperature close to room temperature (Fig. 2C). Detailed explanation of the activation energy calculation is available in SI Appendix, Note 1 and Fig. S3. Hence, the effect of DMSO is to improve the conductivity of a constant number of metallic-like nanocrystals (local order between PEDOT chains) but without delocalizing the electronic wave function of the nanocrystals over the solid. Hence, the electrical transport occurs through percolation paths characterized by a temperature-activated hopping or ultimately, tunneling at high DMSO content (29).

We now turn to the characterization of the HET. For thick (2,000-nm) PEDOT-PSS film with 5% DMSO (dashed line in SI Appendix, Fig. S4), the cyclic voltammogram is composed of a capacitive current (square box i_c) and a faradic current (peak i_f) contribution. The capacitive contribution is proportional to the film thickness, indicating that small ions penetrate through the film, whereas the faradic contribution is constant, with the film thickness suggesting that the redox molecules do not penetrate in the PEDOT-PSS films (SI Appendix, Fig. S5). For thin films of low DMSO content, no visible faradic process takes place despite the presence of Ferro/ferricyanide in the electrolyte due to the low conductivity of the PEDOT-PSS electrode (Fig. 3A and SI Appendix, Fig. S6). The increase of the DMSO content leads to the appearance of faradic peak currents. The clear transition from slow quasiexponential to fast reversible electrode process is indicated by both an increase of peak current and a decrease of the potential of peak-to-peak separation. It is, however, difficult to draw more conclusions based on this dynamic technique due to the complexity of the phenomena and that the electroactivity of the electrode itself depends on its conductivity (SI Appendix, Figs. S7 and S8). To decouple the effect of the capacitive charging of the electrode and the HET kinetics, we measure the faradic current in steady state (Fig. 3B). As mentioned above, since the density of state does not depend on the DMSO content,
we believe that the reading of the faradic current will be directly related to the rate constant of the redox process. The steady-state measurements on films of both high and low conductivities enable the quantification of electrode kinetics through the exchange current (SI Appendix, Note 2, Fig. S9, and Table S1).

We extract the overall rate constant \( k_e \) of the HET and compare with the literature: the \( k_e \) value obtained for the pristine PEDOT-PSS electrode (low conductivity) is five orders of magnitude lower compared with a flat platinum electrode for Ferro/ferricyanide redox couple in aqueous electrolyte (37, 38). However, an increase of DMSO content leads to an increase in the rate constant, and a three-orders of magnitude increase is observed for the PEDOT-PSS electrode with the highest content of DMSO. For inorganic electrodes, there is no obvious correlation between the charge transport within the bulk of the electrode and the rate of HET taking place at the electrode–electrolyte interface; here, we observe a clear linear correlation between the log \( k_e \) and the logarithm of the electrical conductivity of the polymer electrode (Fig. 3C and SI Appendix, Table S1). Our results show that the overall rate of HET \( k_e \) depends on the number of percolations paths, \( P_r \) available for the electronic transport within the polymer electrode that reaches the interface between the electrode surface contact with the electrolyte. Hence, the local rate of electron transfer at the active surface site connected to the percolation path is \( k_{ET} \), and the overall rate is \( k_e = k_{ET}P_r \). This phenomenon is pictured in Fig. 1, showing how percolation paths lead to HET at the polymer–electrolyte interface. That an interfacial HET is governed by the bulk properties of an electrode has never been seen in any other materials than conducting polymers, and we thus can add one more parameter and express the faradic current as

\[
i_f = k_{ET}N_i(E)P_eC_{eq}^{C_s}W^R(E).
\]  

We now apply our findings to TGCs. A TGC (Fig. 4A) is an electrochemical power source that consists of two electrodes in contact with the redox electrolyte solution. The TGC differs from conventional semiconductor-based thermoelectric generators (39), because the key phenomenon generating voltage involves an electron transfer between an electrode and a redox mediator in an electrolyte. The open circuit thermovoltage \( (V_{oc}) \) generated by the application of a temperature difference \( (\Delta T) \) across the TGC can be evaluated by the Seebeck coefficient \( S_e = V_{oc}/\Delta T \) (40, 41), and the entropy change in a redox electrolyte equilibrium is the main origin of the thermovoltage (42, 43). The Ferro/ferricyanide redox couple has been widely studied for TGCs due to its good solubility in water and high exchange currents on metal electrodes, yielding a high Seebeck coefficient \( (1.4 \text{ mV/K}) \) (39, 43–47). In our device, both electrodes are based on PEDOT-PSS without metal underneath. In those electrochemical devices, the electrolyte composition is constant \( 0.4 \text{ M } K_2Fe(CN)_6/K_4Fe(CN)_6 \), and only one parameter is varied: the conductivity of the two PEDOT-PSS electrodes.

Before applying a temperature gradient, we first study the electrochemical phenomena at one PEDOT-PSS electrode in the isothermal condition by impedance spectroscopy. In other words, we replace one of the PEDOT-PSS electrodes with a Pt mesh and characterize on the remaining polymer electrode. The equivalent circuit used to extract the physical parameters has been proposed by others for similar systems (46, 48) (SI Appendix, Figs. S10 and S11 and Table S2). The components are the following: electron transfer resistance \( (R_{ET}) \), Ohmic serial resistance (electrolyte and interfacial contact resistances; \( R_s \)), and bulk electrode resistance \( (R_e) \). Note that \( R_p \) values are similar to the measured \( R_s \) of the polymer film for all values of DMSO content, thus supporting the physical meaning of the circuit component (SI Appendix, Fig. S12). For all DMSO
contents, \( R_{ET} \) is dominant compared with \( R_F \) and \( R_S \) (Fig. 4B). In other words, the electron transfer process is decoupled from the charge transport process in the polymer electrode in all samples, because its kinetics is slower, and it will limit the measured current. Interestingly, the shape of the evolution of \( R_{ET} \) vs. DMSO follows a parallel evolution to \( R_F \), indicating that the two different phenomena are correlated. This fully supports the independent observation that the rate of electron transfer is related to the conductivity of the films (Fig. 3C). \( R_{ET} \) dramatically changes from 100,000 to 100 Ωcm². The kinetics of the dynamic equilibrium for the electron transfer established at the electrode–electrolyte interface is expected to control the TGC performance.

Under temperature gradient between the two PEDOT-PSS electrodes, a thermovoltage is developed, and Seebeck coefficients of 1.43–1.46 mV/K for all PEDOT-PSS blends are extracted, which is in a good agreement with previously reported values with other metal electrodes (39, 49, 50). The independence of the Seebeck coefficient with the DMSO content illustrates the thermodynamic control in the thermovoltage measurements. A dynamic equilibrium of the Ferro/ferricyanide redox process is established even for films of poor conductivity. In contrast, a significant effect of the secondary dopant DMSO content is observed on the power generated by the TGCs acquired at a temperature difference \( \Delta T = 30 \, ^\circ\text{C} \) (Fig. 4C and SI Appendix, Fig. S13). The power increases from 0.01 to 140 mW/m² on DMSO increase, which indicates that the efficiency of the thermogenerator is limited by the electron transfer kinetics. Hence, the four-orders of magnitude increase in both the maximum power density of the TGC and the electrical conductivity of the electrodes (Fig. 24) is correlated to the electron transfer resistance in the impedance data (Fig. 4B) and the exchange current in the Tafel plot (Fig. 3C). The maximum power can be further enhanced by increasing the thickness of the PEDOT-PSS electrode (from a dispersion with 5 wt % DMSO) from 220 to 4,000 nm to reach 410 mW/m², which is like the maximum power obtained with platinum electrodes (435 mW/m²). Note, however, that this increase of power is not due to the penetration of the redox molecule in the film (SI Appendix, Fig. S5) but due to the transition to another limiting mechanism. Indeed, for thin film electrodes (thickness = 220 nm), the limiting phenomenon is the charge transfer resistance, while for thicker films, \( R_S \) becomes limiting (SI Appendix, Fig. S14) as found for PEDOT-Tos [poly(3,4-ethylenedioxythiophene)-tosylate] (S11). Hence, for the same power output, PEDOT electrodes provide a material price/power that is ~120 times lower than the Pt electrode (SI Appendix).

**Conclusion**

In summary, we study the HET between a polymer electrode of semiconducting character and a reactant in solution. While HET at inorganic semiconductor surface strongly depends on the density of state at the Fermi level, we show that a set of semiconducting polymer electrodes with approximately the same surface density of state displays a three-orders of magnitude difference in the electron transfer rate. This reveals that the density of percolation paths for the electronic transport within the polymer electrode dictates the overall rate of electron transfer at the electrolyte–polymer electrode interface. We apply this finding to optimize the power generated by a TGC on heat-to-electricity conversion. The power can be enhanced by four orders of magnitude, demonstrating the importance of the nature of the charge transport within the conducting polymer electrodes. We believe that this finding is of key importance to optimize all of the HET for various emerging electrochemical technologies that could benefit from the unique features of conducting polymers (stability in acid medium, no surface oxide, low temperature and solution processing, porosity, abundant atomic elements). Examples of emerging technologies are fuel cells, dye-sensitized solar cells, and industrial electrochemical synthesis with low-cost conducting polymer electrodes.
Materials and Methods
The PEDOT-PSS film was deposited on the Si substrates by spin coating the commercial aqueous dispersion Clevis PH1000 including 0.1 wt % of silane cross-linker GOPS (details can be found in SI Appendix). The films were about 200-nm thick as measured by atomic force microscopy (SI Appendix). The electrical conductivity was measured by a four-probe method, and the temperature dependence was measured in a vacuum cryogenic probe station (SI Appendix). Cyclic voltammetry and chronopotentiometry were performed in an aqueous solution of 10 mM K₃[Fe(CN)₆]/K₄[Fe(CN)₆] in 1 M KCl as supporting electrolyte against Ag/AgCl reference electrode (details are in SI Appendix). Details on optical absorption and UV–visible spectroscopy as well as TGC apparatus can be found in SI Appendix.

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