Phase Transitions in Operational Risk

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In this paper we explore the functional correlation approach to operational risk. We consider networks with heterogeneous a-priori conditional and unconditional failure probability. In the limit of sparse connectivity, self-consistent expressions for the dynamical evolution of order parameters are obtained. Under equilibrium conditions, expressions for the stationary states are also obtained. The consequences of the analytical theory developed are analyzed using phase diagrams. We find co-existence of operational and non-operational phases, much as in liquid-gas systems. Such systems are susceptible to discontinuous phase transitions from the operational to non-operational phase via catastrophic breakdown. We find this feature to be robust against variation of the microscopic modelling assumptions.

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I. INTRODUCTION

Management and mitigation of risk events are major concerns for banks. The goal is two-fold: first, quantitatively assess the risk in terms of potential financial loss and second, develop solutions to control and buffer the impact of these losses. To facilitate systematic analysis, risk events are broadly classified into three categories: (i) Market Risk (MR), (ii) Credit Risk (CR) and (iii) Operational Risk (OR). MR refers to fluctuations in stock indices, changes in interest rates, foreign exchange parties or commodity (e.g., gold, oil, etc) prices. CR refers to loan defaults when companies go bankrupt. Research on understanding risk and developing sophisticated models has traditionally focused on MR and CR, while OR was initially subsumed under “other”, non credit or market risks. Subsequent spectacular catastrophes including the bankruptcy of the Orange County municipality, California, USA in 1994 [1] and the collapse of Baring Investment Bank, London, United Kingdom in 1995 [2], which were neither attributed to MRs nor CRs, helped establish OR as a risk category of its own. The Basel Committee for Banking Supervision (BCBS), an international regulatory body, now stipulates that banks must explicitly reserve a portion of equity capital against OR.

Note that the above listed three main risk categories are not intended to be exhaustive. Other risk categories exist (e.g., liquidity risk is an important category for bank management, or fiduciary and compliance risks, which arise from the judicial responsibilities of the banks’ customers). Depending on circumstances, these may indeed outweigh the importance of the three ‘main’ risk categories mentioned above.

The Basel II document [3], released by BCBS in 2001 and revised in 2005, is a guideline on banking regulation. Under Basel II, OR is defined as the risk of losses resulting from inadequate or failed internal processes, people and systems, or from external events. From a practical and management perspective it is reasonable to categorize ORs as events based on causes and specific effects. Possible categorizations, as described in [3], are: (i) human processing errors, (ii) human decision errors, (iii) system (software or hardware) errors, (iv) process design errors, (v) fraud and theft and (vi) external damages.

To quantify the capital that must be allocated for operational failures, Basel II suggests three methodologies: (i) Basic Indicator Approach (BIA), (ii) Standardized Approach (SA) and (iii) Advanced Measurement Approach (AMA). Under BIA, the required capital is determined by taking 15% of the banks’ average gross income over the previous three years. The SA is only slightly more advanced in that the banks’ operations are divided into 8 business lines, each of which has a specific weight. The required capital is calculated as the weighted average of the (non-negative) gross income from the business lines over the previous three years. Under AMA, banks are responsible for designing their own measurement approach and setting assumptions on the loss distribution. However, banks must demonstrate that their approach captures potentially severe “tail” loss events. The use of external and internal loss data as well as internal expertise is permitted in the evaluation.

Under AMA, the Loss Distribution Approach (LDA) is a sophisticated and popular measure. Herein the required capital is determined by the Value-at-Risk (VaR) [4] over all OR categories. VaR is defined, over a specified risk horizon \( T \) as the loss not exceeded with probability \( q \), in excess of the expected loss that can happen under normal economic conditions. The VaR is thus dependent on the nature of the loss frequency and severity distributions.

Choices for the loss severity distribution function include log-normal, Gamma, Beta and Weibull distributions. A Poisson or negative binomial distribution is often used for the loss frequency distribution. The common approach to estimate the loss distribution is to first assume that OR categories are independent. Subsequently, for each category \( i \), one draws a realization \( N_i \) from the
loss frequency distribution and samples $N_i$ realizations of the loss severity $X_i^m$ ($m = 1, \ldots, N_i$). The loss is then calculated as

$$L_i = \sum_{m=1}^{N_i} X_i^m.$$  

Finally, drawing a histogram of outcomes of normalized $L_i$, one obtains the loss distribution for each category. The capital to be allocated is the VaR of the sum of losses over all categories, given a time horizon $T$ and confidence level $q$.

Subsequent developments in modeling OR have focused on incorporating more realistic loss frequency and severity distributions. In [6], Chernobai and Rachev provide evidence in favor of using Stable Paretian Distributions. Similarly, in [6], Allen et. al., incorporate a Generalized Pareto Distribution.

A critical point concerning LDA and subsequent analytical development is the assumption of independence between various OR categories. However, a moment’s reflection will lead us to realize that these categories, which may be conceptualized as sets of processes, are functionally dependent.

To put this in context we illustrate it with the collapse of Baring Investment Bank. A trader in the banks’ Singapore office hid trading losses by forging official documents (trading decision errors resulting in fraud), thereby allowing discrepancies to go unnoticed by higher officials. Pleased with his apparent performance, the bank made a mistake of retaining him as Chief Trader (process design error in light of the bank not having adequate checks and balances, which led to decision errors by the banks’ managers). In December 1994, compounded by increasing losses and in order to overcome the predicament, the trader bet on the Tokyo Stock Exchange, assuming that the index would not fall below 19,000. However, in January 1995 a devastating earthquake hit the Japanese city of Kobe, resulting in the index to plummet by 7% in one week, and below 19,000 (series of decision errors compounded by external damage). By the time the bank discovered what had happened, the trader had lost $1.3 billion, effectively bankrupting the bank.

This example shows that the assumption of independence between processes is not realistic. Recent efforts have therefore been directed to capture the influence of interactions between processes. A model based on the lattice gas analogy was proposed in [2]. Investigations were conducted via Monte Carlo (MC) simulations and using a mean-field approximation for a homogeneous and fully connected process network. The model parameters were shown to be related to both conditional and unconditional failure probabilities. Furthermore, avalanches of process failures were shown to be possible through bubble nucleation in first order phase transition. This model was subsequently elaborated by Leipold and Vanini [8]. In [9], Clemente and Ramano present a case study, substantiated by MC simulations, incorporating realistic dependences between processes.

At this point, it is important to note that the concept of OR is not restricted to the banking industry, but is also systemic to any large economy or commerce. Over the past few decades, markets have been subject to considerable de-regulation and globalization. These forces, coupled with an increasing reliance on sophisticated information technology have allowed businesses to develop more efficient operational practices, which include Business Processes Outsourcing (BPO) and automated data collection, storage and retrieval techniques. These advances have led to increased mutual dependencies between economic processes, resulting in a heightened susceptibility for catastrophic breakdowns and thus significant financial losses. A thorough understanding of the dynamics of interacting process networks is more than ever desirable.

Therefore, in consideration of the above, the aim of the present paper is twofold: (i) provide detailed analytical underpinning for the main findings of [2], and (ii) in doing so, broaden the scope of that investigation to highlight the fact that the possibility of catastrophic breakdown in networks of interacting processes is a robust phenomenon under variation of the underlying model assumptions.

The remainder of the paper is organized as follows. In Sec. II we introduce our model. Sec. III begins by introducing three distinct techniques, covering non-equilibrium dynamics and equilibrium statistical mechanics, which we use to derive dynamical evolution and stationary state solutions. In Sec. IIIA we detail a heuristic solution, valid in the case of asymmetric, uncorrelated interactions between processes. In Sec. IIIB a more systematic study, based on a Generating Functional Analysis (GFA) and valid for an arbitrary degree of interaction symmetry is provided. In Sec. IIIC we investigate the stationary behavior of a fully symmetric network using techniques from equilibrium statistical mechanics. In Sec. IV we produce phase diagrams for process networks, exhibiting regions in parameter space where operational phases coexist with non-operational phases. We subsequently evaluate a loss distribution and thereby extend the scope of LDA to account for correlations between processes. Finally, in Sec. V we provide concluding remarks and describe possible extensions for further work. Some of the more technical details from GFA and analysis of stationary states are relegated to appendices A and B respectively.

II. MODEL DEFINITIONS

In this section we describe the statistical model of interacting processes, which was previously introduced in [2]. Each process $i$ ($i = 1, \ldots, N$) is defined by its state at time $t$, which is referred to as $\eta_i(t)$. A two-state model is considered: $\eta_i(t) \in \{0,1\}$. At time $t$ a process can be either up and running, $\eta_i(t) = 0$, or broken down, $\eta_i(t) = 1$. In order to maintain a stable running state over the time increment $t \rightarrow t + 1$, each process $i$ needs
to receive support at time $t$, which is denoted $U_i(t)$ and takes the form

$$U_i(t) = \vartheta_i - \sum_j J_{ij} \eta_j(t) - \xi_i(t).$$  \hfill (2)

Here, $\vartheta_i \in \mathbb{R}$ denotes the average support received by a process in a fully functional environment, while $J_{ij} \in \mathbb{R}$ represents the impact on the average support if process $j$ breaks down. Finally, the $\xi_i(t)$ are zero mean, Gaussian random fluctuations, which represent non-systematic internal and external perturbations to the environment (e.g., fire, earthquake, voltage fluctuations within electric supplies, etc.).

A process breaks down in the next time step if $U_i(t) < 0$. Thus, the dynamics of a processes’ state is given by

$$\eta_i(t+1) = \Theta \left( \sum_j J_{ij} \eta_j(t) - \vartheta_i + \xi_i(t) \right),$$  \hfill (3)

where $\Theta(\ldots)$ represents the Heaviside function. The time-step, is taken to represent one day.

By suitable rescaling of the average support and impact parameters $\vartheta_i$ and $J_{ij}$, respectively, the $\xi_i(t)$ can be taken to have unit variance. The rescaled variables, can then be related to the unconditional and conditional failure probabilities. Consider a situation wherein all processes are working at time-step $t$. The probability that process $i$ will break down in the next step, referred to as $p_i$, is given by integrating Eq. (3) over the noise $\xi_i(t)$.

Thus

$$p_i = \Phi(-\vartheta_i),$$  \hfill (4)

where $\Phi(x) = \frac{1}{2} \left(1 + \text{erf} \left( \frac{x}{\sqrt{2}} \right) \right)$. Similarly, defining $p_{ij}$ as the probability process $i$ will break down in the next step, given that currently process $j$ is broken down while all others are working, gives us

$$p_{ij} = \Phi(J_{ij} - \vartheta_i).$$  \hfill (5)

These relations may be inverted to obtain expressions for the model parameters in terms of the failure probabilities,

$$\vartheta_i = -\Phi^{-1}(p_i), \quad J_{ij} = \Phi^{-1}(p_{ij}) - \Phi^{-1}(p_i).$$  \hfill (6)

In general, this model is not analytically tractable. In investigations, this model is not analytically tractable. In our framework, each process does not interact with all, but instead a fraction of the other processes. We explicitly incorporate this feature by decomposing

$$J_{ij} = c_{ij} \tilde{J}_{ij},$$  \hfill (7)

where $c_{ij} \in \{0,1\}$ are connectivity coefficients and $\tilde{J}_{ij} \in \mathbb{R}$ describes the magnitude of impact. We assume $c_{ij} = c_{ji}$. While this is a reasonable assumption, it is important to realize that the impact magnitudes are not necessarily symmetric. For example, consider a main-frame computer connected to a dummy terminal computer. If the main-frame crashes, we cannot use the terminal. However, if the terminal computer breaks down, it is highly unlikely that it would effect the operations of the main-frame. Thus, in general, $\tilde{J}_{ij} \neq \tilde{J}_{ji}$. The connectivity coefficients are described by the following distribution

$$P(c_{ij}) = \left( 1 - \frac{c}{N} \right) \delta_{[c_{ij},0]} + \frac{c}{N} \delta_{[c_{ij},1]},$$  \hfill (8)

where $c \in \mathbb{R}^+$ is the average connectivity per process. In the case of finite $N$, the adjacency matrix describes an Erdős-Rényi random graph [10].

In what follows we consider the extreme dilution limit, wherein $N \to \infty$, $c \to \infty$, such that $c/N \to 0$. This is achieved by taking $c = \mathcal{O}(\log(N))$. This assumption has important implications on the structure of the connectivity graph. Firstly, each node in the graph will be connected to a vanishing fraction of the total number of nodes. Secondly, the length of a loop is $\mathcal{O}(\log(N))$. Thus, taking $N \to \infty$, the probability of finding loops of finite length tends to 0. The environment about each node is locally tree-like.

The magnitude of the impact parameters $\tilde{J}_{ij}$ are taken to be quenched, i.e., fixed random quantities. To allow for the thermodynamic limit $N \to \infty$ and $c \to \infty$ the mean and variance of the $\tilde{J}_{ij}$ must scale with $c$. We put

$$\tilde{J}_{ij} = \frac{J_0}{c} + \frac{J}{\sqrt{c}} x_{ij},$$  \hfill (9)

where the $x_{ij}$ are zero mean and unit variance random variables. The mean and variance of $\tilde{J}_{ij}$ are parameterized by $J_0 \in \mathbb{R}$ and $J \in \mathbb{R}$ respectively. We note that if $J_0 > 0$ the interactions between processes are, on average, supportive, which is the regime of interest. Secondly small $J$ suppresses the probability of having of negative $\tilde{J}_{ij}$. In addition, small $J$ reduces the effects of frustration [12]. Finally, we choose the $x_{ij}$ to be independent in pairs and have the following moments

$$\bar{x}_{ij} = 0, \quad \bar{x}_{ijkl} = \delta_{[i,k]} \delta_{[j,l]} + \alpha \delta_{[i,l]} \delta_{[j,k]}.$$  \hfill (10)

The parameter $\alpha \in [0,1]$ describes the degree of correlations between $\tilde{J}_{ij}$ and $\tilde{J}_{ji}$, with fully symmetric interactions given by $\alpha = 1$. 


III. MODEL SOLUTIONS

Here we investigate the dynamics and the stationary states of the model introduced in Sec. [11]

For uncorrelated interactions, i.e., $\alpha = 0$, a solution can be obtained by a heuristic argument, following lines of reasoning previously used to study statistical mechanics of disordered systems, in particular neural networks [11]-[13]. This solution relies on assuming weak and negligible correlations between the quenched and dynamic random variables, which cannot be easily justified in a rigorous manner.

A more exact and formal treatment of the dynamics is possible using generating functional analysis (GFA) [14]. This technique facilitates the evaluation of order parameters ab initio and is applicable for arbitrary degree of correlation, $\alpha$, between processes. In addition, it provides a non-trivial check to show that the heuristic solution is exact.

Finally, for fully symmetric networks, $\alpha = 1$, we study the stationary states using techniques from equilibrium statistical mechanics. This requires the use of thermal instead of Gaussian noise.

A. Heuristic Solution

We begin by observing that interactions of process $i$ are described by the local field $h_i(t) = \sum_j c_{ij} x_j \eta_j(t)$. The state $\eta_j(t+1)$ depends on the $\eta_j(t)$ that contribute to $h_i(t)$. Furthermore, the states $\eta_j(t)$ depend on $\eta_j(t-1)$ through $h_j(t-1)$. There is feed-back in the system, which induces correlations between the local fields at time $t$.

In the present mean-field approach we evaluate the statistics of $h_i(t)$ by appealing to the law of large numbers and central limit theorem. The applicability of these tools relies on assuming that the contributions to $h_i(t)$ are weakly correlated if not independent. The condition $\alpha = 0$ together with the limit of sparse connectivity, $c/N \to 0$, entail that contributions to $h_i(t)$ are uncorrelated at finite times. This allows us to describe $h_i(t)$ as a Gaussian random quantity.

Concentrating on the local field $h_i(t)$, we incorporate the definition of $J_{ij}$ given by Eq. (10) to rewrite the expression as

$$ h_i(t) = \frac{J_0}{c} \sum_j c_{ij} \eta_j(t) + \frac{J}{\sqrt{c}} \sum_j c_{ij} x_j \eta_j(t). \quad (11) $$

In the limit $N \to \infty$, the mean and variance of $h_i(t)$ are given by

$$ \langle h_i(t) \rangle_{i,j;x_{ij}} \approx \frac{J_0}{N} \sum_j \langle \eta_j(t) \rangle_{i,j;x_{ij}}, \quad (12) $$

and

$$ \sigma^2(h_i(t)) \approx \frac{J^2}{N} \sum_j \langle \eta_j(t) \rangle_{i,j;x_{ij}}. \quad (13) $$

respectively, [11]. The angled brackets $\langle \ldots \rangle$ refer to the average over noise terms, $\xi_i(t)$, while the over-bar, $\langle \ldots \rangle_{i,j;x_{ij}}$ refers to the average over the coupling parameters, $c_{ij}$ and $x_{ij}$.

In deriving these results, one assumes that correlations between dynamical degrees of freedom $\eta_k(t)$ and the $x_j$ and $c_{ij}$ that characterize disorder, are negligible. This allows factoring of averages of the form $c_{ij} x_{ij} \langle \eta_k(t) \rangle_{i,j;x_{ij}} \approx c_{ij} x_{ij} \langle \eta_k(t) \rangle_{i,j;x_{ij}}$. The fraction $m(t+1)$ of failed processes at the next time-step, $t + 1$, evaluates to

$$ m(t+1) = \frac{1}{N} \sum_j \eta_j(t+1) = \frac{1}{N} \sum_j \Theta(h_j(t) - \vartheta_j + \xi_j(t)). \quad (14) $$

In the limit $N \to \infty$, Eq. (14) can be evaluated by appealing to the law of large numbers as a sum of averages over the random variables $h_j(t)$, $\xi_j(t)$ and $\vartheta_j$. As $h_j(t)$ and $\xi_j(t)$ are Gaussian random variables, their sum is also Gaussian, with mean $J_0 m(t)$ and variance $1 + J^2 m(t)$. Thus, first performing the joint average over local fields and noise, we obtain

$$ m(t+1) = \frac{1}{N} \sum_i \Phi \left( \frac{J_0 m(t) - \vartheta_i}{\sqrt{1 + J^2 m(t)}} \right). \quad (15) $$

Finally, we note that the only $i$ dependence in Eq. (15) comes from the $\vartheta_i$. Thus, in the limit $N \to \infty$, we obtain an average over the $\vartheta$ distribution [22], which we denote by $\langle \ldots \rangle^\vartheta$, giving

$$ m(t+1) = \Phi \left( \frac{J_0 m(t) - \vartheta}{\sqrt{1 + J^2 m(t)}} \right)^\vartheta. \quad (16) $$

We have thus obtained a simple closed expression for the evolution of the fraction of failed processes in the network.

B. Generating Functional Analysis

Investigations of the dynamical properties of a system are conducted systematically employing GFA [14], which provides tools for the evaluation of correlation and response functions in terms of a characteristic functional of path-probabilities. Performing the average over bond disorder in the sum over dynamical trajectories, one obtains a family of effective single site processes parameterized by $\vartheta$,

$$ \eta(t+1) = \Theta \left( J_0 m(t) + \alpha J^2 \sum_{s < t} G(t,s) \eta(s) \right. $$

$$ - \vartheta + \phi(t) + h(t). \quad (17) $$
These single site processes exhibit memory, via the response function, \( G(t, s) \), and are driven by colored Gaussian noise, \( \{ \phi(t) \} \), self consistently determined via
\[
\langle \phi(s) \phi(t) \rangle = \delta_{|s-t|} + J^2 q(s, t), \tag{18}
\]
\[
m(t) = \langle \eta(t) \rangle^\beta, \tag{19}
\]
\[
q(s, t) = \langle \eta(s) \eta(t) \rangle^\beta, \tag{20}
\]
\[
G(t, s) = \frac{\partial m(t)}{\partial h(s)}, \quad s < t. \tag{21}
\]
Here \( \langle (\ldots) \rangle \) refers to average over \( \{ \phi(t) \} \). The external field \( h(t) \) is primarily introduced to define the response function via Eq. (21). The derivation and interpretation of these parameters are provided in appendix A. The order parameter \( m(t) \), equation (19), describes the dynamics of the fraction of failed processes.

In the case \( \alpha = 0 \), we recover Eq. (16), which was derived via purely heuristic reasoning in Sec. IIIA. However, for arbitrary \( \alpha \), the response function complicates averaging over \( \{ \phi(t) \} \). In these cases, numerical results for Eqs. (19)-(21) are obtained using Eissfeller-Oppe (EO) simulations [15]. A key concern in performing the simulations is producing colored noise. Our method uses a Cholesky decomposition [16] of the noise covariance matrix. We construct colored noise as a linear combination of white noises weighted with elements of the Cholesky matrix.

**C. Equilibrium Statistical Mechanics**

A process network of the form described in Sec. III is guaranteed to achieve a stationary probability distribution for its microscopic states, if we assume \( \alpha = 1 \) and use thermal noise, distributed according to
\[
p(\xi) = \frac{1}{\sqrt{2\pi}} \exp\left(-\frac{\beta^2}{4}\xi^2\right), \tag{22}
\]
instead of Gaussian noise, as in Eq. (9). The parameter \( \beta \) is called the inverse temperature. The Gibbs-Boltzmann equilibrium distribution is characterized by the Hamiltonian,
\[
H(\eta) = -\sum_{i < j} c_{ij} \tilde{J}_{ij} \eta_i \eta_j + \sum_i \tilde{\theta}_i \eta_i. \tag{23}
\]
A solution of this model, with disorder variables \( \tilde{\eta} \), \( c_{ij} \) and \( \tilde{x}_{ij} \), given by Eqs. (11)-(13), is achieved with techniques used to solve the Sherrington-Kirkpatrick (SK) spin-glass [17] model. In this spirit, the Replica-Symmetric (RS) order parameters are obtained from the extensive free energy as solutions of the following pair of self-consistency equations,
\[
m = \int Dz \Phi^\beta(h_{RS}^\beta), \tag{24}
\]
\[
q = \int Dz \Phi^2(h_{RS}^\beta), \tag{25}
\]
where,
\[
\Phi^\beta(x) = \frac{1}{2} \left( 1 + \tanh\left( \frac{\beta x}{2} \right) \right), \tag{26}
\]
and
\[
h_{RS} = -\tilde{\theta} + J_0 m + J \sqrt{q} z + \frac{\beta J^2}{2} (m - q). \tag{27}
\]
Derivations are provided in appendix B. The order parameter \( m \), Eq. (24), describes the stationary fraction of failed processes.

The stability of this solution against Replica Symmetry Breaking (RSB) was checked by verifying that the Hessian at the RS saddle point is positive definite [18]. This requires that the so called replicon eigenvalue, given by
\[
\lambda = 1 - (\beta J)^2 \int Dz \left( \Phi^\beta(h_{RS}) - \Phi^\beta(h_{RS}^2) \right)^2 \tag{28}
\]
is positive. We found that for networks with parameter settings as investigated in section IV below, i.e., a-priori homogeneous failure probability, \( p = 0.01 \), RSB occurs only for \( J \gtrsim 4.630 \). Hence, the regime of interest with small frustration, say \( J < 0.5 \), is well within the stable domain.

Comparison between the equilibrium \( m \), Eqs. (24)-(27), and the long-term stationary behavior of \( m(t) \), Eq. (19), with \( \alpha = 1 \) is possible, once we scale the thermal noise appropriately to match properties of the Gaussian noise. In principle, such a matching may be accomplished in several different manners, each of which leads to similar results. Here, we prescribe that the thermal noise should have the same unit variance as the Gaussian noise used in the microscopic dynamics. This leads to \( \lambda = \pi/\sqrt{3} \).

**IV. RESULTS**

In this section we explore the consequences of theory developed above, with particular emphasis on phase diagrams of resulting stationary states.

For simplicity, we assume homogeneous \( \tilde{\theta} \), which allows us to drop the corresponding \( \tilde{\theta} \)-averages. We justify this simplification by noting that any other non-trivial \( \tilde{\theta} \)-distribution, for example Gaussian with mean \( \overline{\tilde{\theta}} \) and variance \( \sigma^2_{\tilde{\theta}} \), would not alter the qualitative behavior of the observables. We illustrate this for the fully asymmetric network, \( \alpha = 0 \). To evaluate Eq. (14) we exploit that the sum of Gaussian terms, \( h_j(t) - \tilde{\theta}_j + \xi_j(t) \), is itself a Gaussian with mean \( J_0 m(t) - \overline{\tilde{\theta}} \) and variance \( 1 + J^2 m(t) + \sigma^2_{\tilde{\theta}} \). This gives
\[
m(t+1) = \Phi \left( \frac{J_0 m(t) - \overline{\tilde{\theta}}}{\sqrt{1 + J^2 m(t) + \sigma^2_{\tilde{\theta}}}} \right). \tag{29}
\]
Up to a shift in the variance of the noise, this equation describes the evolution of a system of asymmetrically coupled processes with uniform \( \vartheta_i = \mathcal{J} \).

For stationary states for the fully symmetric network, \( \alpha = 1 \), obtained from equilibrium statistical mechanics, a similar structural shift can be shown to apply.

A. Fully Asymmetric Network

Assuming long-term stationary behavior for the observable \( m(t) \), Eq. (10), we drop the time index. The curve in the first panel of Fig. 1 depicts \( m \) as a function of \( J_0 \). We note the coexistence of a low-\( m \) operational and a high-\( m \) non-operational phase for \( J_0 \) values bounded by the lower and upper critical values, \( J_0^c \approx 3.971 \) and \( J_0^c \approx 14.814 \), respectively. Within this interval there is an unstable branch of solutions, with intermediate \( m \)-values, represented by the back-bending part of the curve.

The behavior is parameterized by the unconditional failure probability \( p \). This dependence is made explicit in the phase diagram, shown in the second panel of Fig. 1. The upper curve marks the \( J_0^c \)-boundary of the low-\( m \) operational phase, while the lower curve represents the \( J_0^c \)-bound of the high-\( m \) non-operational phase. A transition between the two phases is discontinuous for \( p < p_c \approx 0.102 \) and becomes continuous exactly at \( p_c \). Beyond \( p_c \), the phases lose their separate identity, much as in a liquid-gas system. In the region of small \( J \), which is of interest for OR, say \( J < 0.5 \), the behavior shown in Fig. 1 is fairly insensitive to variation in \( J \).

B. Partially Symmetric Network

Investigations of partially symmetric networks were conducted via: (i) evaluating the first three time-steps of the effective single site dynamics exactly, starting with random initial conditions, given by \( m(0) \), and (ii) conducting EO simulations of the dynamics to higher time-steps.

The results presented in Fig. 2 were produced for an intermediate value of the symmetry parameter \( \alpha = 0.5 \).

The exact evaluation of the first three time-steps is depicted in the first panel. The curves represent the fraction of failed processes at each time as a function of \( J_0 \). We observe that these curves intersect the line \( m(0) = 0.3 \) at exactly the same point, \( J_0^c \approx 6.001 \), which marks a change in behavior. For \( J_0 < J_0^c \), a convergence to the low-\( m \) operational phase is realized, while for \( J_0 > J_0^c \), we observe convergence to the high-\( m \) non-operational phase. Except in the immediate vicinity of \( J_0^c \), the convergence is rapid and achieved within the first three time-steps.

To investigate the behavior beyond the first three time-steps, we used EO simulations to propagate Eqs. (17)-(21) up to \( t = 30 \). This time horizon is sufficient to achieve stationarity, except in an infinitesimal neighborhood of the critical \( J_0^c \) values. In the second panel of Fig. 2 we provide the stationary \( m \) value as a function of \( J_0 \). The coexistence of the low-\( m \) operational phase and high-\( m \) non-operational phase is observed for intermediate \( J_0 \), in complete analogy to the fully asymmetric case. These \( J_0 \) values are bounded between the critical values, \( J_0^c \approx 3.966 \) and \( J_0^c \approx 14.513 \). Here, the unstable branch, represented by the back-bending curve, is computed by locating the \( J_0^c \) values that separate the regions of convergence towards either the operational or non-operational phase. It is noteworthy that we obtain an almost identical unstable branch if we use the exact forms of the first three time-steps.
C. Fully Symmetric Network

Investigations of fully symmetric networks, \( \alpha = 1 \), were conducted via: (i) appealing to equilibrium statistical mechanics, Eqs. (22) - (27), and (ii) EO simulations.

To compare the two results, we scaled the thermal noise to have unit variance, resulting in \( \beta = \pi/\sqrt{3} \).

In the first panel of Fig. 3 we plot values of \( m \) as a function of \( J_0 \). The dotted line was produced using techniques from equilibrium statistical mechanics. The solid line was produced from EO simulations for \( \alpha = 1 \). In both cases, the coexistence of a low-\( m \) operational and a high-\( m \) non-operational phase, for intermediate values of \( J_0 \), is observed. The coexistence regions are given by the intervals \([4.113, 20.478]\) and \([3.966, 14.512]\) for the networks with thermal and Gaussian noise, respectively. Whereas the critical values for the non-operational phase are almost identical, a significant discrepancy for the upper critical value is observed. This discrepancy must be attributed to differences in the nature of the noises.

In the second panel of Fig. 3 we produce the phase diagram for the system with thermal noise. We observe a similar qualitative behavior to the fully asymmetric case, \( \alpha = 0 \), given in Fig. 1. The critical point \( p_c \approx 0.110 \) is very close to that from the asymmetric case, \( p_c \approx 0.102 \).

D. Loss Distribution and Capital Requirement

Finally, in Fig. 4 we show a loss distribution for a network of 50 processes. The stationary macroscopic variable \( m \) represents the fraction of failed processes. Alternatively, \( m \) represents the probability that a process breaks down within an interacting environment. That is, effects of functional correlations with other processes in the network have been accounted. We proceed by draw-
ing from a binomial distribution, for each process $i$ the number $k_i$ of days in a year the process fails. The daily failure probability is given by $m$. Failures on different days are taken to be independent. Subsequently binning the sum of $k_i$ realizations from the log-normal loss severity distribution for each process, we obtain the loss distribution for the network.

In the thermodynamic limit one would expect the loss distribution to be a Gaussian by the central limit theorem. However, for $N = 50$ the distribution is still heavily skewed to the right with a “fat” tail and rare realizations of extremely large losses.

A qualitative study was also conducted to determine the effects of functional correlations on the VaR. Recall that under LDA, VaR is used to determine the capital to be allocated for ORs. We consider a fully asymmetric network with homogeneous unconditional and conditional failure probabilities. First, for a non-interacting environment, the loss frequency distribution is binomial with daily failure probability, $m = \phi(\theta)$. Using the construction as described afore, the loss distribution was evaluated and VaR$_0$ calculated for a confidence level $q = 99.99\%$. Next, for the interacting environment, the loss frequency distribution is once more binomial. However the daily failure probability, $m$ is now given by the stationary operational solution to Eq (10), as a function of $J_0$. The loss distributions are evaluated and the VaR determined for the same confidence level. In Fig. (5) we plot the ratio VaR to VaR$_0$, as a function of $J_0$. We consider a network of 50 processes with $J = 0.2$, over a risk horizon of 365 days. The data is fitted against a quadratic curve, represents by the dashed line.

V. CONCLUSION

In this paper we investigated networks of interacting and functionally correlated processes as models of OR in large organizations. In the limit of sparse connectivity, we studied the dynamical evolution and stationary states of these networks using analytical methods and numerical simulations. A heuristic approach was used to exactly solve the dynamics of fully asymmetric networks. GFA and EO simulations were used to study the dynamics of partially and fully symmetric networks. Finally, stationary states for fully symmetric networks were analyzed using techniques from equilibrium statistical mechanics.

The initial variant of our model was formulated on a sparsely connected random graph, for which the average connectivity per process $c$, satisfies, $c/N \rightarrow 0$ in the thermodynamic limit. However, it should be noted that the assumption needed to carry through the GFA as described in this paper is $c \gg 1$ rather than sparseness. Hence, the analysis applies to networks with non-sparse connectivity with $c = O(N)$ as well.

From all investigated cases, we found that there exists a range of interactions, characterized by $J_0$, for which an operational phase coexists with a non-operational phase. This raises the possibility of spontaneous catastrophic breakdown in such process networks. We demonstrated this feature is robust and invariant under a broad range of changes in microscopic details.

Phase diagrams were produced for fully asymmetric networks using the exact dynamical theory and for fully symmetric networks using the results from equilibrium statistical mechanics. Similar behavior is observed in both cases. For sufficiently small values of a-priori failure probability $p$, the operational and non-operational phases are distinct; a discontinuous transition between phases can be induced, for instance, by changing $J_0$. There is
a critical failure probability \( p_c \), where the transition is continuous. We also note that the two critical values, \( p_c \approx 0.102 \) for the fully asymmetric case, and \( p_c \approx 0.110 \) for the fully symmetric case are very close.

Given parameter values \( p \), \( J_0 \) and \( J \) for which stable operational and non-operational phases coexist, one can further qualify these phases as either absolutely stable or meta-stable. These notions are well known for thermodynamic equilibrium systems. Therein, the absolutely stable phase is the one with the lower value of the free energy. Parameter values at which absolutely stable phases become meta-stable (or vice-versa) are the locations of proper first-order equilibrium phase transitions.

Of the systems studied in this paper, however, only the fully symmetric network with thermal noise can be characterized as a thermodynamic equilibrium system. Nevertheless, the notion of absolute stability and meta-stability of coexisting dynamic stationary states can be carried over to the other systems in the following manner. One denotes by \( \tau_N \) the fraction of the total time a system of size \( N \) spends in one of the coexisting stationary states, either operational or non-operational. A state is characterized as absolutely stable if \( \lim_{N \to \infty} \tau_N = 1 \), and as meta-stable, if \( \lim_{N \to \infty} \tau_N = 0 \). A meta-stable phase can nevertheless be dynamically stable in the sense that, once in such a phase, a system will typically stay in that phase for a very long time \( t_N \) which diverges as \( N \to \infty \).

These concepts have a significant implication for OR, that play out for large (but finite) process networks that exhibit dynamically and absolutely stable operational phases, coexisting with a meta-stable non-operational phase. As there are no detectable precursors for transitions between them, the assessment must be performed by stress tests, wherein artificial strains and fluctuations are introduced into the system and the effects are measured. The present model readily lends itself to tests of this type.

In our investigation, we only considered Poisson random graphs with large average connectivity \( J \). As a consequence, there is little heterogeneity in the local environment about each process. However, it is known that realistic process networks exhibit a richer topology, which includes the presence of “hubs”. In such situations, the connectivity scales according to a power-law and the analytical treatment would follow lines of reasoning used in small-world networks \([21]\). A similar qualitative behavior with coexistence of operational and non-operational phases is nevertheless expected.

**APPENDIX A: GENERATING FUNCTIONAL ANALYSIS**

In this appendix we use Generating Functional Analysis (GFA) \([14]\) to formally solve the model dynamics and systematically justify the heuristic solution in Sec. \[III A\].

We begin by introducing the generating functional over source fields, \( \psi \),

\[
Z[\psi] = \langle \exp\left(-i \sum_{i=1}^{N} \sum_{t=0}^{T} \psi_i(t) \eta_i(t)\right) \rangle. \tag{A1}
\]

The angled brackets refer to the average over all “paths”, which are trajectories of microscopic states. Explicitly,

\[
Z[\psi] = \sum_{\{\eta(t)\}} P[\{\eta(t)\}] \exp\left(-i \sum_{t=0}^{T} \sum_{i=1}^{N} \psi_i(t) \eta_i(t)\right) \tag{A2}
\]

where \( P[\{\eta(t)\}] = P[\eta(0), \ldots, \eta(T)] \) denotes the probability of all paths over the risk horizon, \( T \). The generating functional can be used to compute expectation values and correlation functions as

\[
\langle \eta_i(t) \rangle = i \frac{\partial Z[\psi]}{\partial \psi_i(0)}|_{\psi \equiv 0}, \tag{A3}
\]

\[
\langle \eta_j(s) \eta_i(t) \rangle = i^2 \frac{\partial Z[\psi]}{\partial \psi_j(s) \partial \psi_i(t)}|_{\psi \equiv 0}. \tag{A4}
\]
We are interested in evaluating the generating functional for a typical realization of disorder. This is achieved by averaging Eq. (A2) over \( c_{ij} \) and \( x_{ij} \). To proceed, we exploit that the path probability measure has a Markov structure.

\[
P(\{\eta(t)\}) = P(\eta(0)) \prod_{t=0}^{T-1} P(\eta(t+1) | \eta(t)), \quad (A5)
\]

with transition probabilities

\[
P(\eta(t+1) | \eta(t)) = \prod_{i=1}^{N} \int \mathcal{D}\xi_i(t) \delta[\eta_i(t+1), f_i(t)]. \quad (A6)
\]

The \( \xi_i(t) \) are unit mean and zero variance Gaussian random variables. Furthermore,

\[
f_i(t) = \Theta \left( \sum_j c_{ij} \hat{J}_{ij} \eta_j(t) + h_i(t) - \vartheta_i + \xi_i(t) \right). \quad (A7)
\]

1. Average Over Fast Noise

We conduct the \( \xi_i(t) \) integral and disorder average by first extracting these contributions \( u_i(t) = \sum_j c_{ij} \hat{J}_{ij} \eta_j(t) + \xi_i(t) \) from the Heaviside function in \( f_i(t) \). This is facilitated utilizing a Dirac \( \delta \)-function and its normalization property to give us

\[
Z[\psi] = \sum_{\eta(0)} \cdots \sum_{\eta(T)} P(\eta(0)) \prod_{t=0}^{T-1} \prod_{i=1}^{N} \left[ \int \frac{d\tilde{u}_i(t)}{2\pi} \exp \left( -i \tilde{u}_i(t) \left( u_i(t) - \sum_j c_{ij} \hat{J}_{ij} \eta_j(t) \right) \right) \right. \\

\left. - \left[ \hat{u}_i(t)^2 \right] \delta[\eta_i(t+1), f_i(t)] \right] \exp \left( -i \sum_{t=0}^{T} \sum_{i=1}^{N} \psi_i(t) \eta_i(t) \right). \quad (A8)
\]

In equation (A8), we redefined \( f_i(t) \) to denote,

\[
f_i(t) = \Theta \left( u_i(t) + h_i(t) - \vartheta_i \right). \quad (A9)
\]

The “hatted” conjugate terms, \( \tilde{u}_i(t) \), are a consequence of a Fourier representation of the \( \delta \)-function.

2. Disorder Average

The disorder average, which factorize in pairs, affects \( c_{ij} \) and \( x_{ij} \). We localize the terms involved in the following definition

\[
\mathcal{D}^{c_{ij}, x_{ij}} = \prod \exp \left( \sum_{t=0}^{T-1} \sum_{i,j} c_{ij} \left( \hat{u}_i(t) J_{ij} \eta_j(t) \right) \right. \\

\left. + i \hat{u}_j(t) J_{ji} \eta_i(t) \right) \]. \quad (A10)

We first perform the \( c_{ij} \) average, and proceed by taking the Taylor expansion of the exponential in the limit \( c \gg 1 \). Next, taking the \( x_{ij} \) average, re-exponentiating and keeping the dominant terms, we obtain

\[
\mathcal{D}^{c_{ij}, x_{ij}} = \exp \left( N \left[ J_0 \sum_{t=0}^{T-1} k(t) m(t) \right. \right. \\

\left. \left. + \frac{J^2}{2} \sum_{s,t=0}^{T-1} \left( Q(s,t) q(s,t) \right. \right. \right. \\

\left. + \alpha G(t,s) G(t,s) \right] \right). \quad (A11)
\]

which depends on a set of macroscopic variables defined as

\[
m(t) = \frac{1}{N} \sum_{i=1}^{N} \eta_i(t), \quad k(t) = \frac{1}{N} \sum_{i=1}^{N} i \tilde{u}_i(t), \quad q(s,t) = \frac{1}{N} \sum_{i=1}^{N} \eta_i(s) \eta_i(t),
\]

\[
Q(s,t) = \frac{1}{N} \sum_{i=1}^{N} i \tilde{u}_i(s) i \tilde{u}_i(t), \quad G(t,s) = \frac{1}{N} \sum_{i=1}^{N} \tilde{u}_i(s) \eta_i(t).
\]

To achieve site factorization in the evaluation of \( Z[\psi]^{c_{ij}, x_{ij}} \), one proceeds as usual by enforcing the above
definitions of the macroscopic variables via Dirac δ-function identities and their Fourier representations. This subsequently generates a set of conjugate variables to the above set of macroscopic variables. The averaged generating functional is expressed in the following compact form

\[ \mathcal{Z}[\psi] = \int \mathcal{D}\{\psi\} \exp \left\{ N [\Xi_1 + \Xi_2 + \Xi_3] \right\} . \]  

(A12)

Here, \( \mathcal{D}\{\psi\} \) represents taking the integral over all the macroscopic order parameters, and their conjugates. The functions \( \Xi_1, \Xi_2 \) and \( \Xi_3 \), appearing in the exponential of Eq. (A12), are defined as

\[ \begin{align*}
\Xi_1 &= J_0 \sum_{t=0}^{T-1} k(t) m(t) + J_2^2 \sum_{s,t=0}^{T-1} \left( Q(s,t) q(s,t) \right. \\
&\quad + 1 \left. \sum_{s,t=0}^{T-1} \left( q(s,t) \tilde{q}(s,t) + Q(s,t) \tilde{Q}(s,t) \right) \right) + G(t,s) \tilde{G}(t,s), \\
\Xi_2 &= \frac{1}{N} \sum_{i=1}^{N} \log \left\{ \sum_{\{\eta(0)\}} P(\eta(0)) \right\} \\
\Xi_3 &= \frac{1}{N} \sum_{i=1}^{N} \left\{ \int \frac{d\mu(t) d\nu(t)}{2\pi} \delta_{[\eta(t+1), f_i(t)]} \right\} \\
&\quad \exp \left( -\mathcal{S} - \frac{1}{N} \sum_{i=1}^{N} \psi_i(t) \eta(t) \right) \right\}. \\
\end{align*} \]

(A13)

(A14)

(A15)

Here \( \mathcal{S} \) denotes the dynamic action,

\[ \begin{align*}
\mathcal{S} &= \sum_{t=0}^{T-1} \left( -\hat{u}_i(t) u_i(t) - \hat{u}_i(t)^2/2 - \hat{m}(t) \eta_i(t) \right) \\
&\quad + \hat{Q}(s,t) \tilde{Q}(s,t) \tilde{Q}(s,t) \tilde{Q}(s,t) \tilde{Q}(s,t) \eta_i(s) \eta_i(s) \right) .
\end{align*} \]

Thus, we have transformed the generating functional into an integral with leading order in the exponential of \( N \), which can be evaluated using the saddle point technique. The contribution \( \Xi_3 \), Eq. (A15) describes an ensemble of independent dynamical processes.

3. Saddle Point Equations

At the saddle point, the macroscopic observables of interest resolve to

\[ \begin{align*}
m(t) &= \frac{1}{N} \sum_{i=1}^{N} \langle \eta(t) \rangle_{(i)} , \\
q(s,t) &= \frac{1}{N} \sum_{i=1}^{N} \langle \eta(s) \eta(t) \rangle_{(i)} \\
G(t,s) &= \frac{1}{N} \sum_{i=1}^{N} \langle \hat{u}_s \eta(s) \rangle_{(i)} .
\end{align*} \]

(A16)

The \( \langle (\ldots) \rangle_{(i)} \) represent averages over the dynamics of effective single site processes and have the form

\[ \langle (\ldots) \rangle_{(i)} = \frac{\sum_{\{\eta(t)\}} P(\eta(0)) \left( \prod_{t=0}^{T-1} \left[ \int \frac{d\mu(t) d\nu(t)}{2\pi} \delta_{[\eta(t+1), f_i(t)]} \right] \right) (\ldots) e^{-\mathcal{S}}}{\sum_{\{\eta(t)\}} P(\eta(0)) \left( \prod_{t=0}^{T-1} \left[ \int \frac{d\mu(t) d\nu(t)}{2\pi} \delta_{[\eta(t+1), f_i(t)]} \right] \right) e^{-\mathcal{S}}} \]

(A17)

Averages involving a conjugate field \( \hat{u}_i(t) \) describe response functions, i.e., perturbations of expectations with respect to an external field. Thus, averages involving only conjugate fields describe a perturbation of a constant and are set to zero. Furthermore, by causality, \( G(t,s) \), which describes the response of the fraction of failed processes at time \( t \), to a perturbation at time \( s \), vanishes \( \forall s \geq t \). Using the saddle point identities for other order parameters and their conjugates, we obtain \( \Xi_1 = 0 \) and \( \Xi_2 = 0 \) at the saddle point. The dynamic action reduces to

\[ \begin{align*}
\mathcal{S} &= \sum_{t=0}^{T-1} \left( \hat{u}_i(t) \left[ u(t) - J_0 m(t) \right] \right. \\
&\quad - \alpha J^2 \sum_{s,t} G(t,s) \eta(s) \left. \right) \\
&\quad + \frac{1}{2} \sum_{s,t=0}^{T-1} \hat{m}(t) \tilde{m}(s) \left( J^2 q(s,t) + \delta_{[s,t]} \right) .
\end{align*} \]

(A18)
and it corresponds to single site processes, with dynamics
\[ \eta(t + 1) = \Theta \left( J_0 m(t) + \alpha J^2 \sum_{s < t} G(t, s) \eta(s) - \vartheta + \phi(t) + h(t) \right). \]

(A19)

We note the following: (i) There is now a dependence on the fraction \( m(t) \) of failed processes. (ii) If there is a degree of symmetry, \( \alpha \neq 0 \), the dynamics is non-Markovian, with memory given by the response function \( G(t, s) \). (iii) The noise \( \phi(t) \) is colored, with \( \langle \phi(t) \rangle = 0 \) and \( \langle \phi(s) \phi(t) \rangle = J^2 q(s, t) + \delta_{|s-t|} \).

We recall that the only \( i \) dependence in Eqs. (A16)-(A17) comes from \( \vartheta_i \) in \( f_i(t) \), Eq. (A9). By the law of large numbers the empirical averages in Eq. (A16) can be evaluated as an average over the \( \vartheta \) distribution,
\[ \lim_{N \to \infty} \frac{1}{N} \sum_i \langle (\ldots) \rangle_i = \int d\vartheta p(\vartheta) \langle (\ldots) \rangle = \langle (\ldots) \rangle. \]

The saddle point results, Eq. (A16) thus take the form
\[ m(t) = \langle (\eta(t))^\vartheta \rangle, \]
\[ q(s, t) = \langle \eta(s) \eta(t) \rangle^\vartheta, \]
\[ G(t, s) = \frac{\partial m(t)}{\partial h(s)}. \]

(A20) (A21) (A22)

In the case \( \alpha = 0 \), there is no memory effect. Eq. (A20) is by itself sufficient to describe the dynamics. Evaluating the \( \phi(t) \) average we get
\[ m(t + 1) = \Phi \left( J_0 m(t) - \vartheta \sqrt{1 + J^2 m(t)} \right). \]

(A23)

However, for \( \alpha \neq 0 \), memory \( G(t, s) \) is relevant. The non-linear nature of the evolution Eqs. (A19)-(A20) precludes a simple analytical characterization of the long-time asymptotic stationary states.

**APPENDIX B: EQUILIBRIUM STATISTICAL MECHANICS**

For fully symmetric networks, \( \alpha = 1 \), and thermal noise, \( \xi \), distributed as
\[ p(\xi) = \frac{1}{2} \beta \text{sech}^2 \left( \frac{\beta \xi}{2} \right), \]

(B1)
equilibrium statistical mechanics can be employed to study the stationary behavior. The resulting Gibbs-Boltzmann equilibrium distribution for observing a given microscopic state is characterized by the Hamiltonian
\[ H(\eta) = - \sum_{i < j} c_{ij} J_{ij} \eta_i \eta_j + \sum_i \vartheta_i \eta_i. \]

(B2)

We obtain the stationary macroscopic order parameter, \( m \), describing the fraction of failed processes from the free energy per process. In the limit, \( N \to \infty \), the free energy is expected to be self averaging over the disorder,
\[ f = \lim_{N \to \infty} - \frac{T}{N} \log \left( \sum_{\eta} \exp \left( N \left( \frac{\beta J_0}{2} \sum_{\mu} m^2_{\mu} + \frac{\beta^2 J^2}{4} \sum_{\mu, \nu} q^2_{\mu, \nu} \right) - \beta \sum_{\mu} \sum_i \vartheta_i \eta^\mu_i \right) \right). \]

(B3)

The above quenched average is calculated using the “replica trick”,
\[ f = \lim_{N \to \infty} \lim_{n \to 0} - \frac{T}{N} \log \left( \sum_{\eta} \exp \left( N \left( \frac{\beta J_0}{2} \sum_{\mu} m^2_{\mu} + \frac{\beta^2 J^2}{4} \sum_{\mu, \nu} q^2_{\mu, \nu} \right) - \beta \sum_{\mu} \sum_i \vartheta_i \eta^\mu_i \right) \right). \]

(B4)

In a manner similar to that presented in appendix A2, we take the average over \( c_{ij} \) and \( x_{ij} \). The resulting expression for the free energy is,
\[ f = \lim_{n \to 0} \lim_{N \to \infty} - \frac{T}{n N} \log \sum_{(\eta^\mu)} \exp \left( N \left( \frac{\beta J_0}{2} \sum_{\mu} m^2_{\mu} + \frac{\beta^2 J^2}{4} \sum_{\mu, \nu} q^2_{\mu, \nu} \right) - \beta \sum_{\mu} \sum_i \vartheta_i \eta^\mu_i \right), \]

(B5)

where,
\[ m_{\mu} = \frac{1}{N} \sum_{i=1}^N \eta^\mu_i, \quad q_{\mu, \nu} = \frac{1}{N} \sum_{i=1}^N \eta^\mu_i \eta^\nu_i. \]

(B6)

The indices \( 1 \leq \mu, \nu \leq n \) label the replicas. The above parameters were introduced to achieve site factorization, in a manner identical to that in appendix A2. Subsequently, the free energy is expressed as an integral, which is evaluated via the saddle point technique.
\[
f = \lim_{n \to 0} \lim_{N \to \infty} -\frac{T}{N} \log \int \mathcal{D}(\ldots) \exp \left( \sum_{\alpha} m_{\alpha}^2 + \frac{\beta^2 J^2}{4} \sum_{\alpha, \epsilon} q_{\alpha \epsilon}^2 + 1 \sum_{\alpha} \hat{m}_\alpha m_\alpha + 1 \sum_{\alpha, \epsilon} \hat{q}_{\alpha \epsilon} q_{\alpha \epsilon} \right) \quad \text{(B7)}
\]

Here, the effective Hamiltonian is given by
\[
-\beta H_{\text{eff}} = -\beta \vartheta \sum_{\alpha} \eta^\alpha - i \sum_{\alpha} \tilde{m}_\alpha \eta^\alpha
- i \sum_{\alpha, \epsilon} \tilde{q}_{\alpha \epsilon} \eta^\alpha \eta^\epsilon .
\quad \text{(B8)}
\]

1. Saddle Point Equations

We obtain the following expressions for the order parameters,
\[
m_\alpha = \langle \eta^\alpha \rangle, \quad q_{\alpha \epsilon} = \langle \eta^\alpha \eta^\epsilon \rangle ,
\quad \text{(B9)}
\]

where,
\[
\langle \ldots \rangle = \frac{\sum_{\eta^\alpha} \exp (-\beta H_{\text{eff}})}{\sum_{\eta^\alpha} \exp (-\beta H_{\text{eff}})} .
\quad \text{(B10)}
\]

A solution of the order parameters is achieved with the following replica symmetric ansatz,
\[
m_\alpha = m, \forall \alpha \quad \text{(B11)}
q_{\alpha \epsilon} = m \delta_{[\alpha, \epsilon]} + q (1 - \delta_{[\alpha, \epsilon]}), \forall \alpha, \epsilon
\quad \text{(B12)}
\]

On application of the ansatz and Gaussian linearization of quadratic replica terms, we obtain
\[
m = \int Dz \Phi_\vartheta (h_{\text{RS}}) ,
q = \int Dz \Phi_\vartheta^2 (h_{\text{RS}}) ,
\quad \text{(B13, B14)}
\]

where, \( \Phi_\vartheta (x) = \frac{1}{2} \left( 1 + \tanh \left( \frac{\beta z}{2} \right) \right) \) and
\[
h_{\text{RS}} = -\vartheta + J_0 m + J \sqrt{q} z + \frac{\beta J^2}{2} (m - q) .
\quad \text{(B15)}
\]

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