Structure-property relations of co-doped bismuth layer-structured Bi$_{3.25}$La$_{0.75}$(Ti$_{1-x}$Mo$_x$)$_3$O$_{12}$ ceramics

Pasinee Siriprapa$^1$, Anucha Watcharapasorn$^{1,2}$ and Sukanda Jiansirisomboon$^{1,2}$*

Abstract

In this work, the fabrication and investigation of substituting higher-valence Mo$^{6+}$ for Ti$^{4+}$ ion on the B-site of La$^{3+}$-doped Bi$_4$Ti$_3$O$_{12}$ [BLT] structure to form Bi$_{3.25}$La$_{0.75}$(Ti$_{1-x}$Mo$_x$)$_3$O$_{12}$ [BLTM] (when $x = 0, 0.01, 0.03, 0.05, 0.07, 0.09$, and 0.10) ceramics were carried out. X-ray diffraction patterns of BLTM ceramics indicated an orthorhombic structure with lattice distortion, especially with a higher concentration of a MoO$_3$ dopant. Microstructural investigation showed that all ceramics composed mainly of plate-like grains. An increase in MoO$_3$ doping content increased the length and thickness of the grain but reduced the density of the ceramics. Electrical conductivity was found to decrease, while the dielectric constant increased with Mo$^{6+}$ doping concentration. Ferroelectric properties were found to be improved with increasing MoO$_3$ content and were optimized at $x = 0.1$.

Keywords: ceramics, X-ray diffraction, dielectric properties, microstructure, ferroelectricity

Background

In recent years, the family of bismuth layer-structured ferroelectrics has received much attention as the candidate for ferroelectric random access memories. An extensively studied bismuth titanate (Bi$_4$Ti$_3$O$_{12}$ [BIT]) is a member of the Aurivillius family that can be represented by a general formula (Bi$_2$O$_3$)$_{2}$[A$_{m-1}$(B)$_m$O$_{3m+1}$] which consists of (Bi$_2$O$_2$)$^{2+}$ sheets alternating with (Bi$_2$Ti$_3$O$_{10}$)$^{2+}$ perovskite-like layers [1]. BIT has large spontaneous polarization along the $a$ axis (approximately 50 μC/cm$^2$), low processing temperature, high Curie temperature, and is a Pb-free material [2]. However, it still has high leakage current and domain pinning due to defects, such as Bi vacancies accompanied by oxygen vacancies [3,4]. To overcome these problems, A-site substitution by a replacement of volatile Bi with rare earth or other metal oxide additives is often necessary for ferroelectric property improvement. For example, ions of La [5], V [6], Nd [7], and Pr [8] have been used to substitute the Bi ion in a BIT bulk material or thin film without destroying its layered structure. Bu et al. [9] prepared La-doped BIT thin films by pulsed laser deposition and reported that these films were appropriate for non-volatile random access memory devices because of their high remanent polarization and low leakage current. In recent years, SimÕes et al. [10] reported that a doping content of $x = 0.75$ in Bi$_4$La$_x$Ti$_3$O$_{12}$ [BLT] showed an improvement of the fatigue endurance upon a repeated cyclic electric field which emphasized its possible use in FRAM applications. However, these BLT ceramics still showed a rather high leakage current. Attention has thus been paid to investigate this material in order to overcome this disadvantage. Wang et al. [11] reported that substitution of high-valence ions, such as Mo$^{6+}$, in BLT thin films or known as Mo$^{6+}$ co-doped BLT thin films led to polarization fatigue improvement, high remanent polarization, and lower leakage current density.

In terms of ceramics, there has been no detailed study on a new co-doped bismuth-layered structure based on Mo$^{6+}$-doped Bi$_{3.25}$La$_{0.75}$Ti$_3$O$_{12}$. In this study, therefore, effects of MoO$_3$ doping concentration on the phase, microstructure, and electrical properties (i.e., conductivity, dielectric, and ferroelectric properties) of BLT ceramics produced by a conventional solid-state mixed-oxide method are reported and discussed.

Methods

A perovskite bismuth-layered structure based on Bi$_{3.25}$La$_{0.75}$(Ti$_{1-x}$Mo$_x$)$_3$O$_{12}$ [BLTM] ($x = 0, 0.01, 0.03,$...
0.05, 0.07, 0.09, and 0.10) powders was prepared using a solid-state mixed-oxide method. Starting binary oxide powders, i.e., Bi$_2$O$_3$ (> 98%, Fluka; Sigma-Aldrich Corporation, St. Louis, MO, USA), La$_2$O$_3$ (99.98%, Fluka), TiO$_2$ (> 99%, Riedel-de Haën; Sigma-Aldrich Corporation, St. Louis, MO, USA), and MoO$_3$ (99.9%, Fluka), were ball-milled and calcined at 750°C for 4 h. The slurry was transferred to a spherical flask and placed in a shell freezer. The flask was rotated in an ethanol bath for at least 1 h. The flask of frozen slurry was then immediately transferred to a vacuum dryer and dried for 24 h. After the ice was sublimated, fine dried powder was produced. The BLTM powders were then pressed under a uniaxial hydraulic pressure of 5.5 MPa with a few drops of 3 wt.% polyvinyl alcohol used as a binder. The pressed samples were sintered at temperatures in a range of 1,000 to 1,150°C for 4 h. Optimum sintering temperature for producing highest-density ceramics was determined, and the samples were selected for further characterization. Phases of selected ceramics were characterized using an X-ray diffractometer [XRD] (X-pert, PANalytical B.V., Almelo, The Netherlands) with CuK$_\alpha$ radiation. Density was measured by Archimedes' method. The ceramics were polished and thermally etched at a temperature of 150°C below the optimum sintering temperature for 15 min dwell time prior to microstructural investigation using a scanning electron microscope [SEM] (JEOL JSM-6335F, JEOL Ltd., Akishima, Tokyo, Japan). Average grain size was determined using a mean linear intercept method from SEM micrographs. Electrical conductivity measurement was done at 1 kHz using an LCZ meter. Dielectric constant [$\varepsilon_r$] and loss tangent [tanδ] were measured at room temperature with a frequency between 1 to 100 kHz using LCR Hitester 3532-50 (Hioki, Ueda, Nagano, Japan). Ferroelectric hysteresis polarization-electric field [P-E] loops were determined using a computer-controlled modified Sawyer-Tower circuit. Remanent polarization [$P_r$], maximum polarization [$P_{max}$], coercive field [$E_c$], maximum field [$E_{max}$], and loop squareness [$R_{sq}$] values were evaluated from the loops.

**Results and discussion**

X-ray diffraction patterns of BLTM calcined powders with various MoO$_3$ doping concentration are shown in Figure 1. The patterns were well matched with the ICSD No. 150091 database of Bi$_{3.25}$La$_{0.75}$Ti$_3$O$_{12}$, which indicated an existence of a single orthorhombic phase without a detectable second phase. The densification data of BLTM ceramics sintered at various temperatures between 1,000 to 1,150°C (not shown here) indicated that the optimum sintering condition for achieving maximum density was 1,100°C and 4 h dwell time. These ceramic samples were selected for phase characterization whose results are shown in Figure 2. All XRD patterns basically indicated an orthorhombic structure similar to those found in the powders. It could be seen from the XRD patterns of samples with high MoO$_3$ doping content (≥0.05 mol) that X-ray intensities of a particular set of planes, i.e., (004), (006), (008), (0012), (0014), and (0016), were relatively higher than those in the samples with lower MoO$_3$ content. This suggested that the preferred orientation of BLT crystallites was increased due to Mo$^{6+}$ ions.

Density and grain size measured in terms of grain length are presented in Table 1. Increasing MoO$_3$ doping content was found to reduce density values. Typical morphologies of polished and thermally etched surfaces of BLTM ceramics are shown in Figure 3. All samples showed plate-like morphology with variation in grain size and orientation. It was obvious that both grain length and thickness gradually increased with increasing MoO$_3$ content. It was also recognized that MoO$_3$ played an important role in increasing anisotropies of the surface and grain boundary energies and mobilities [12,13]. The sample containing MoO$_3$ ≥ 0.05 mol showed grains which tended to stack and align in the same direction. For those containing a smaller amount of MoO$_3$ (≤0.03 mol), their grains became smaller with a more randomized grain orientation. This observation corresponded well to the XRD result (Figure 1b) where the preferred orientation of (001)-type planes occurred in BLTM ceramics with high MoO$_3$ concentration.

The electrical conductivity of BLT and BLTM ceramics measured at frequencies of 1, 50, and 100 kHz is shown in Figure 4. The conductivity values of BLTM were lower than those of undoped BLT ceramic regardless of measurement frequency. This seemed to be the effect of the donor dopant on the electrical conductivity of a p-type material in which electrons were introduced when Mo$^{6+}$ substituted for Ti$^{4+}$ site in BLT ceramics. For BLT, electrical conductivity at low frequencies, i.e., 1 and 50 kHz, was much lower than that at 100 kHz. This could be explained by the contributions of space-charge polarization at the grain boundary and ionic polarization caused by ionic motion [15-17]. It has been reported that a low-frequency conductivity response was dominated by the grain boundary, while a high-frequency conductivity response was dominated by ferroelectric crystalline layers inside the grains. Hence, it could be speculated that the low-frequency dependence was dominated by an interacting charge carrier system, which was similar to an ionic conductor. For BLTM ceramics, the apparent decrease in frequency dependence of alternating current...
[AC] conductivity seemed to correspond well to an increase in grain size which in turn reduced the amount of grain boundaries. Among BLTM samples whose average grain size was about the same, these ceramics therefore showed similar frequency dependence and magnitude of AC conductivity.

Room temperature $\varepsilon_r$ and dielectric tan$\delta$ of the BLTM ceramics are shown in Figure 5, and the values are listed in Table 1. It was experimentally observed that the dielectric constant gradually increased from 127 for BLT to a maximum of 414 for 0.03 mol MoO$_3$ doping content. An increase in dielectric constant of BLTM was expected to be caused by partial substitution of Mo$^{6+}$ ion for Ti$^{4+}$ ions in the B-site which implied that the MoO$_3$ dopant reduced space charges or ionic conduction and resulted in a low leakage current which, in turn, resulted in an enhanced dielectric constant. A further increase in MoO$_3$ content over 0.03 mol was found to decrease the dielectric constant, possibly due to overcompensation of charge which caused a slight increase in electrical conductivity.

Ferroelectric P-E hysteresis loops of MoO$_3$-doped BLT ceramics were measured at room temperature with an applied field of 60 kV/cm and a frequency of 50 Hz as shown in Figure 6. For undoped BLT, no hysteresis behavior was observed. It was previously reported by Zhang et al. [17] that an electric field as high as 530 kV/cm was required to pole a BLT thin film at room temperature. In addition, Kim et al. [18] showed that hysteresis behavior was observed when the polarization was carried out using an applied field of approximately 60 kV/cm at 50°C. Therefore, it seemed that a relatively high leakage current as seen from its high electrical conduction was the main reason for the non-hysteresis behavior observed in the
BLT ceramic. Remanent polarization and coercive field were found to increase after Mo$^{6+}$ was doped into BLT ceramics which were similar to a previous investigation on Nb-doped BLT ceramics [18]. Because of the temperature and field dependence of ferroelectric properties, i.e., remanent polarization and coercive field of ceramics [19,20], these parameters were normalized in a form of $P_r/P_{\text{max}}$ and $E_c/E_{\text{max}}$ values, where $P_{\text{max}}$ was the polarization value.

Table 1 Physical, dielectric, and ferroelectric properties of Bi$_{3.25}$La$_{0.75}$(Ti$_{1-x}$Mo$_x$)$_3$O$_{12}$ ceramics

| Mo content (mol) | Density (g·cm$^{-3}$) | Grain length (µm) | Dielectric properties$^a$ | Ferroelectric properties |
|------------------|-----------------------|-------------------|--------------------------|-------------------------|
|                  |                       |                   | $\varepsilon_r$ | tanδ | $P_r/P_{\text{max}}$ | $E_c/E_{\text{max}}$ | $R_{sq}$ |
| 0                | 7.46 ± 0.03           | 2.59 ± 0.61       | 127          | 0.001 | 0.05             | 0.05             | 0.06     |
| 0.01             | 7.21 ± 0.01           | 2.66 ± 0.79       | 124          | 0.007 | 0.23             | 0.23             | 0.25     |
| 0.03             | 7.16 ± 0.03           | 4.61 ± 2.13       | 414          | 0.551 | 0.96             | 0.69             | 1.15     |
| 0.05             | 6.74 ± 0.02           | 5.94 ± 2.11       | 402          | 0.519 | 0.74             | 0.58             | 0.84     |
| 0.07             | 6.94 ± 0.02           | 7.73 ± 3.13       | 387          | 0.510 | 0.78             | 0.70             | 1.04     |
| 0.09             | 6.79 ± 0.01           | 8.04 ± 3.31       | 310          | 0.461 | 0.81             | 0.71             | 1.06     |
| 0.10             | 6.83 ± 0.02           | 8.13 ± 2.97       | 308          | 0.349 | 0.96             | 0.57             | 1.15     |

$^a$Measurement was carried out at room temperature and at a frequency 1 kHz. $\varepsilon_r$, dielectric constant; tanδ, loss tangent; $P_r$, remanent polarization; $P_{\text{max}}$, maximum polarization; $E_c$, coercive field; $E_{\text{max}}$, maximum field; $R_{sq}$, loop squareness.
at the $E_{\text{max}}$. These ratios are listed in Table 1 along with the values of $R_{\text{sq}}$. It seemed that Mo$^{6+}$ substituting for Ti$^{4+}$ in the BLT ceramic efficiently decreased the concentration of $V_\circ$, which weakened the influence of domain pinning on polarization [21]. A decrease of space charge ($V_\circ$) density as well as electrical conductivity thus led to the large remanent polarization observed. In addition, changes in domain structures may be one of the factors affecting the coercive field [14]. As can be seen from the figure, the BLTM ceramic with a Mo doping concentration of 0.1 mol showed a well defined P-E hysteresis loop and better squareness than BLT and other BLTM ceramics. This research thus suggested that the partial substitution of Mo$^{6+}$ for Ti$^{4+}$ ions at the B-site could improve ferroelectric properties of BLT ceramics with an optimum ferroelectric behavior obtained at $x = 0.1$ mol.
Conclusions
A new system of co-doped bismuth titanate-layered structure ceramics, i.e., $\text{Bi}_{3.25}\text{La}_{0.75}(\text{Ti}_{1-x}\text{Mo}_x)_3\text{O}_{12}$ or BLTM ($x = 0, 0.01, 0.03, 0.05, 0.07, 0.09, \text{ and } 0.10$), was successfully prepared by a solid-state mixed-oxide method. X-ray diffraction analysis indicated that the MoO$_3$ dopant induced a preferred orientation of the BLT ceramics with changes in lattice constant of its

Figure 4 Relationship between electrical conductivity and MoO$_3$ content of $\text{Bi}_{3.25}\text{La}_{0.75}(\text{Ti}_{1-x}\text{Mo}_x)_3\text{O}_{12}$ ceramics. The measurement was done at frequencies of 1, 50, and 100 kHz.

Figure 5 Relationship between dielectric constant and MoO$_3$ content of $\text{Bi}_{3.25}\text{La}_{0.75}(\text{Ti}_{1-x}\text{Mo}_x)_3\text{O}_{12}$ ceramics. The measurement was done at frequencies of 1, 50, and 100 kHz. The inset shows dielectric loss tangent as a function of MoO$_3$ content.
orthorhombic structure. Grain size was found to increase with increasing MoO₃ doping content. Electrical conductivity of BLTM was slightly decreased by Mo⁶⁺ donor doping due to the charge compensation between the provided excess electrons and inherent holes. This led, consequently, to an increased dielectric constant from 127 for BLT to a maximum of 414 for 0.03 mol MoO₃ doping content. Addition of MoO₃ content also increased the remanent polarization of over 10 μC/cm², and this value was shown to be much larger than that of the BLT ceramic, due to a reduction of domain pinning oxygen vacancies and leakage current loss with optimum ferroelectric properties obtained for the composition of x = 0.1 mol. It is suggested that molybdenum-doped BLT could be an alternative material for potential applications in electronic industries which require a lead-free material having large remanent polarization and lower processing temperatures.

Acknowledgements
This work is financially supported by the Thailand Research Fund (TRF) and the Nation Research University Project under Thailand’s Office of the Higher Education Commission (OHEC). PS would like to thank the Commission on Higher Education for their support through a grant fund under the program Strategic Scholarships for Frontier Research Network for the Ph.D. Program Thai Doctoral degree, Faculty of Science and Graduate School, Chiang Mai University for this research.

Author details
1Department of Physics and Materials Science, Faculty of Science, Chiang Mai University, Chiang Mai, 50200, Thailand 2Materials Science Research Center, Faculty of Science, Chiang Mai University, Chiang Mai, 50200, Thailand

Authors’ contributions
PS carried out the BLTM ferroelectric ceramic experiments, analysis, and writing of the manuscript. AW and SJ participated in the conception and design of the study and revised the manuscript for important intellectual content. All authors read and approved the final version of the manuscript.

Competing interests
The authors declare that they have no competing interests.

Received: 6 September 2011 Accepted: 5 January 2012 Published: 5 January 2012

References
1. Aurivillius B: Mixed bismuth oxides with layer lattices: II. structure of Bi₄Ti₃O₁₂. Arkiv Kemi 1949, 1:499-512.
2. Yao YY, Song CH, Bao P, Su D, Lu XM, Zhu JS, Wang YN: Doping effect on the dielectric property in bismuth titanate. J Appl Phys 2004, 95:3126-3130.
3. Lohkamper R, Neumann H, Arlt G: Internal bias in acceptor-doped BaTiO₃ ceramics: numerical evaluation of increase and decrease. J Appl Phys 1990, 68:4220-4227.
4. Stewart WC, Cosentino LS: Some optical and electrical switching characteristics of a lead zirconate titanate ferroelectric ceramic. Ferroelectrics 1970, 1:149-167.
5. Park BH, Kang BS, Bu SD, Noh TW, Lee JH, Jo WK: Lanthanum-substituted bismuth titanate for use in non-volatile memories. Nature 1999, 401:682-684.
6. Noguchi Y, Miyayama M: Large remanent polarization of vanadium-doped Bi₂Ti₃O₁₂. Appl Phys Lett 2001, 78:1903-1905.
7. Kojima T, Sakai T, Watanabe T, Funakubo H, Saito K, Osada M: Large remanent polarization of (Bi, Nd)₂Ti₃O₁₂ epitaxial thin films grown by metalorganic chemical vapor deposition. Appl Phys Lett 2002, 80:2476-2478.
8. Chen M, Liu ZL, Wang YJ, Wang CC, Yang XS, Yao KL: Ferroelectric properties of PrₓO₁₁-doped Bi₂Ti₃O₁₂. Solid State Commun 2004, 130:735-739.
9. Bu SD, Kang BS, Park BH, Noh TW: Composition dependence of the ferroelectric properties of lanthanum-modified bismuth titanate thin films grown by using pulsed-laser deposition. J Korean Phys Soc 2000, 36: L9-L12.

10. Simões AZ, Quinelato C, Ries A, Stoianovic BD, Longo E, Varela JA: Preparation of lanthanum doped Bi4Ti3O12 ceramics by the polymeric precursor method. Mater Chem Phys 2006, 98:481-485.

11. Wang X, Ishitoh A: Polarization enhancement and coercive field reduction in W- and Mo-doped Bi3.35La0.75Ti3O12 thin film. Appl Phys Lett 2003, 82:2479-2481.

12. Horn DS, Messing GL: Anisotropic grain growth in TiO2-doped alumina. Mater Sci Eng A 1995, 195:196-205.

13. Powers JD, Gaesser AM: In Sintering Technology. Edited by: German RM, Messing GL, Cornwall RG. Marcel Dekker Press, New York; 1996.

14. Zhang L, Chu R, Zhao S, Li Q, Yin Q: Microstructure and electrical properties of niobium doped Bi2Ti2O7 layer-structured piezoelectric ceramic. Mater Sci Eng B 2005, 116:99-103.

15. Shulman HS, Damjanovic D, Setzer N: Niobium doping and dielectric anomalies in bismuth titanate. J Am Ceram Soc 2000, 83:528-532.

16. Kim JS: Effects of Nb on dielectric and ferroelectric properties of bismuth titanate ceramics. Integr Ferroelectr 2006, 79:139-145.

17. Zhang ST, Yuan GL, Wang J, Chen YF, Cheng GJ, Liu ZG: Temperature-dependent effect of oxygen vacancy on polarization switching of ferroelectric Bi2.25La0.75Ti3O12 thin films. Solid State Commun 2004, 132:315-318.

18. Kim JS, Ahn CW, Lee HJ, Kim IW, Jin BM: Nb doping effects on ferroelectric and electrical properties of ferroelectric Bi1.25La0.75Ti1-xMoxO3 ceramics. Ceram Int 2004, 30:1459.

19. Ye Z, Cui MH, Cheng CP, Zhou YC, Zheng XJ, Hu ZS: Simulation of polarization and butterfly hysteresis loops in bismuth layer-structured ferroelectric thin films. J Appl Phys 2006, 100:094101-1-094101-5.

20. Burkhanov AL, Shilnikov AV, Sopit AV, Luchaninov AG: Dielectric and electromechanical properties of (1 - x)PMN-xPZT ferroelectric ceramics. Phys Solid State 2000, 42:936-943.

21. Baudry L: Theoretical investigation of the influence of space charges on ferroelectric properties of PbZrTiO3 thin film capacitor. J Appl Phys 1999, 86:1096-1105.

doi:10.1186/1556-276X-7-42

Cite this article as: Siriprapa et al.: Structure-property relations of co-doped bismuth layer-structured Bi1.25La0.75(Ti1-xMo-x)3O12 ceramics. Nanoscale Research Letters 2012 7:42.

Submit your manuscript to a SpringerOpen journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Immediate publication on acceptance
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at ► springeropen.com