\textbf{\(\gamma\gamma\) physics with the KLOE experiment}

The KLOE-2 Collaboration

F. Archilli,\textsuperscript{a,o} D. Babusci,\textsuperscript{f} D. Badoni,\textsuperscript{n,o} I. Balwierz,\textsuperscript{e} G. Bencivenni,\textsuperscript{f} C. Bini,\textsuperscript{l,m} C. Bloise,\textsuperscript{f} V. Boci,\textsuperscript{m} F. Bossi,\textsuperscript{f} P. Branchini,\textsuperscript{q} A. Budano,\textsuperscript{p,q} S. A. Bulychjev,\textsuperscript{g} P. Campana,\textsuperscript{f} G. Capon,\textsuperscript{f} F. Ceradini,\textsuperscript{p,q} P. Ciambrone,\textsuperscript{f} E. Czerwiński,\textsuperscript{f} E. Dané,\textsuperscript{f} E. De Lucia,\textsuperscript{f} G. De Robertis,\textsuperscript{b} A. De Santis,\textsuperscript{l,m} G. De Zorzi,\textsuperscript{l,m} A. Di Domenico,\textsuperscript{l,m} C. Di Donato,\textsuperscript{h,i} D. Domenici,\textsuperscript{f} O. Erriquez,\textsuperscript{a,b} G. Fanizzi,\textsuperscript{a,b} G. Felici,\textsuperscript{f} S. Fiore,\textsuperscript{l,m} P. Franzini,\textsuperscript{l,m} P. Gauzzi,\textsuperscript{l,m} S. Giovannella,\textsuperscript{f} F. Gonnella,\textsuperscript{n,o} E. Graziani,\textsuperscript{f} F. Happacher,\textsuperscript{f} B. Höistad,\textsuperscript{a} E. Iarocci,\textsuperscript{j,f} M. Jacewicz,\textsuperscript{s} T. Johansson,\textsuperscript{s} V. Kulikov,\textsuperscript{g} A. Kupsc,\textsuperscript{s} J. Lee-Franzini,\textsuperscript{f,r} F. Loddo,\textsuperscript{b} M. Martemianov,\textsuperscript{g} M. Martini,\textsuperscript{f,k} M. Matsuuk,\textsuperscript{c} R. Messi,\textsuperscript{n,o} S. Miscetti,\textsuperscript{f} G. Morello,\textsuperscript{f} D. Moriccianni,\textsuperscript{a} P. Moskal,\textsuperscript{e} F. Nguyen,\textsuperscript{p,q} A. Passeri,\textsuperscript{g} V. Patera,\textsuperscript{j,f} I. Prado Longhi,\textsuperscript{p,q} A. Ranieri,\textsuperscript{b} P. Santangelo,\textsuperscript{f} I. Sarra,\textsuperscript{f} M. Sciubba,\textsuperscript{p,q} B. Sciascia,\textsuperscript{f} A. Sciubba,\textsuperscript{j,f} M. Silarski,\textsuperscript{a} C. Taccini,\textsuperscript{p,q} L. Tortora,\textsuperscript{g} G. Venanzoni,\textsuperscript{f} R. Versaci,\textsuperscript{j,u} W. Wiślicki,\textsuperscript{t} M. Wolke,\textsuperscript{s} J. Zdebik,\textsuperscript{e}

\textsuperscript{a}Dipartimento di Fisica dell’Università di Bari, Bari, Italy.
\textsuperscript{b}INFN Sezione di Bari, Bari, Italy.
\textsuperscript{c}Dipartimento di Fisica dell’Università della Calabria, Cosenza, Italy.
\textsuperscript{d}INFN Gruppo collegato di Cosenza, Cosenza, Italy.
\textsuperscript{e}Institute of Physics, Jagiellonian University, Cracow, Poland.
\textsuperscript{f}Laboratori Nazionali di Frascati dell’INFN, Frascati, Italy.
\textsuperscript{g}Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia.
\textsuperscript{h}Dipartimento di Fisica dell’Università “Federico II”, Napoli, Italy.
\textsuperscript{i}INFN Sezione di Napoli, Napoli, Italy.
\textsuperscript{j}Dipartimento di Scienze di Base ed Applicate per l’Ingegneria dell’Università “Sapienza”, Roma, Italy.
\textsuperscript{k}Dipartimento di Scienze e Tecnologie applicate, Università “Guglielmo Marconi”, Roma, Italy.
\textsuperscript{l}Dipartimento di Fisica dell’Università “Sapienza”, Roma, Italy.
\textsuperscript{m}INFN Sezione di Roma, Roma, Italy.
\textsuperscript{n}Dipartimento di Fisica dell’Università “Tor Vergata”, Roma, Italy.
\textsuperscript{o}INFN Sezione di Roma Tor Vergata, Roma, Italy.
\textsuperscript{p}Dipartimento di Fisica dell’Università “Roma Tre”, Roma, Italy.
\textsuperscript{q}INFN Sezione di Roma Tre, Roma, Italy.
\textsuperscript{r}Physics Department, State University of New York at Stony Brook, USA.
\textsuperscript{s}Department of Nuclear and Particle Physics, Uppsala University, Uppsala, Sweden.
\textsuperscript{t}A. Soltan Institute for Nuclear Studies, Warsaw, Poland.
1. Introduction

The term "γγ physics" stands for the study of the reaction of order α^4 e^+e^- → e^+e^-γ^*γ^* → e^+e^-X, where X is some arbitrary final state allowed by conservation laws. In particular, hadronic states with J^P C = 0^±+ 2^±+ are directly produced through the γ^*γ^* → X subprocess. For quasi-real photons the number of produced events can be estimated from the expression

\[ N = L_{ee} \int dW_{γγ} \frac{dL}{dW_{γγ}} \sigma(γγ → X) \]  

where \( L_{ee} \) is the e^+e^- luminosity, \( W_{γγ} \) the photon-photon center of mass energy (\( W_{γγ} = M_X \)), \( dL/dW_{γγ} \) the photon-photon flux and \( σ \) the cross section into a given final state. By neglecting single powers of \( \ln(E/m_e) \), when the scattered leptons are undetected in the final state one has

\[ \frac{dL}{dW_{γγ}} = \frac{1}{W_{γγ}} \left( \frac{α}{2π} \ln \frac{E}{m_e} \right)^2 f(z) \quad z = \frac{W_{γγ}}{2E} \]  

with \( f(z) \) (Low function) given by

\[ f(z) = (2 + z^2)^2 \ln \frac{1}{z} - (1 - z^2)(3 + z^2) \]  

The \( γγ \) processes studied with KLOE are e^+e^- → e^+e^-γ and e^+e^- → e^+e^-π^0π^0. Past experiments measuring the \( γγ \) cross section for light mesons production, took data at a center of mass energy \( 7 < \sqrt{s} < 35 \) GeV. A low energy e^+e^- factory, such as DAΦNE, compensates the small cross section value with the high luminosity.

DAΦNE is an e^+e^- collider designed to operate at the center of mass energy \( \sqrt{s} \simeq 1.02 \) GeV, namely the φ meson mass. The sample used for the present analyses consists of data taken at \( \sqrt{s} = 1 \) GeV, which allows the reduction of the background from φ decays, with an integrated luminosity of 240 pb^-1.

The KLOE detector consists of a large (3.3 m length and 2 m radius) volume drift chamber, surrounded by a lead-scintillating fibers calorimeter. The detector is inserted in a superconducting coil producing an axial field \( B = 0.52 \) T. Charged particle momenta are reconstructed with resolution \( σ_p/p \simeq 0.4\% \) (\( σ_p/p \simeq 1\% \)) for large (small) angle tracks. Calorimeter energy clusters are reconstructed with energy and time resolution of \( σ_E/E = 5.7\%/\sqrt{E(\text{GeV})} \) and \( σ_t = 57 \) ps/\( \sqrt{E(\text{GeV})} \) ± 100 ps.

Data are processed with a dedicated filter asking for at least two prompt photons, clusters not associated to tracks and propagating with \( |p_t - c| \sim \sim 0 \), with energy \( E > 15 \) MeV and polar angle \( 20^° < θ < 160^° \), the most energetic with \( E > 50 \) MeV, the fraction of energy carried by photons \( R = \langle Σ_e E_e \rangle/\langle E_{\text{cal}} \rangle > 0.3 \) and the total energy in the calorimeter \( 100 < E_{\text{cal}} < 900 \) MeV. The latter requirement rejects low energy background and the high rate processes e^+e^- → e^+e^-, e^+e^- → γγ.
2. Observation of $\gamma\gamma \rightarrow \eta$, with $\eta \rightarrow \pi^+\pi^-\pi^0$

The selection of these events asks for two photons compatible with a $\pi^0$ decay and two tracks with opposite curvature coming from the collision point. The charged pion mass is assigned to the two tracks and a least squares function based on Lagrange multipliers imposes that $\pi^+\pi^-\pi^0$ come from an $\eta$ decay. Therefore most background events are suppressed, except for the irreducible process $e^+e^- \rightarrow \eta\gamma \rightarrow \pi^+\pi^-\pi^0\gamma$, with the monochromatic photon, $E_\gamma = 350$ MeV, lost in the beam pipe. For these background events, however, the correlation between the squared missing mass, $M_{\text{miss}}^2$, and the longitudinal momentum, $p_L$, of the $\pi^+\pi^-\pi^0$ system is rather different than for the signal. Further criteria are applied to suppress processes with photons and $e^+e^-$ in the final state. The Monte Carlo, MC, generator used for the signal allows the full phase space generation of $e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-\eta$ events. Fig. 1 (left) shows the $M_{\text{miss}}^2$ distribution for data fitted with the superposition of MC shapes for signal and background. An independent fit is performed with the $\pi^+\pi^-\pi^0$ longitudinal momentum. Both fits show the same yields for the background processes and more than 600 signal events. The cross section of the irreducible process $e^+e^- \rightarrow \eta\gamma \rightarrow \pi^+\pi^-\pi^0\gamma$ is evaluated using the same sample of data and asking for three photons in the final state. A kinematic fit is performed, requiring energy and momentum conservation, and the improved variables are used to fit in the distribution of the recoil photon energy (Fig. 1 right). The preliminary result is $\sigma_{e^+e^- \rightarrow \eta\gamma}(1 \text{ GeV}) = 0.866 (9)_{\text{stat}} (93)_{\text{syst}} \text{ nb}$, where the systematic uncertainty is given by residual $e^+e^- \rightarrow \omega\pi^0 \rightarrow \pi^+\pi^-\pi^0\pi^0$ background.

The cross section $\sigma_{\gamma\gamma \rightarrow \eta \rightarrow \pi^+\pi^-\pi^0}(1 \text{ GeV})$ is under evaluation.

\begin{figure}[h]
\centering
\includegraphics[width=0.45\textwidth]{figure1a.pdf}
\includegraphics[width=0.45\textwidth]{figure1b.pdf}
\caption{Left: fit of the $M_{\text{miss}}^2$ distribution for the $e^+e^- \rightarrow e^+e^-\eta$ analysis. Main contributions are: $e^+e^-\gamma$ (green) at negative $M_{\text{miss}}^2$ values due to the pion mass assigned to $e^+e^-$ tracks, $\eta\gamma$ (red), signal events (light blue). Right: fit of the recoil photon energy spectrum for the $e^+e^- \rightarrow \eta\gamma$ analysis. The $\eta\gamma$ peak (blue) and the $\omega\gamma$ peak (light blue) are visible; the broad distribution (green) is due to $\omega\pi^0$ events.}
\end{figure}

3. Observation of $\gamma\gamma \rightarrow \eta$, with $\eta \rightarrow \pi^0\pi^0\pi^0$

The main backgrounds for this analysis are annihilation processes with at least four prompt photons in the final state, $e^+e^- \rightarrow \eta\gamma$, $K_S K_L$, $\omega\pi^0$, for which accidental or split calorimeter clusters increase the photon multiplicity. Events with six and only six prompt photons in the final state and with no tracks in the drift chamber are selected. The photons are paired choosing...
the combination which minimizes the $\chi^2$-like variable

$$\chi^2_{6\gamma} = \frac{(m_{\pi^0} - m_{ij})^2}{\sigma_{ij}^2} + \frac{(m_{\pi^0} - m_{mn})^2}{\sigma_{mn}^2} + \frac{(m_{\pi^0} - m_{kl})^2}{\sigma_{kl}^2},$$

(4)

where $m_{ij}$ is the invariant mass of each pair of photons and $\sigma_{ij}$ the resolution. Then, a kinematic fit is performed asking for the six photons invariant mass to be equal to the mass of the $\eta$ meson. A cut is then applied on the energy of the most energetic photon, to reject $e^+e^- \rightarrow \eta\gamma$ events in which the monochromatic photon is detected. The squared missing mass, $M_{\text{miss}}^2$, distribution is well described by the $\gamma\gamma \rightarrow \eta \rightarrow \pi^0\pi^0\pi^0$ and $\eta(\rightarrow \pi^0\pi^0\pi^0)\gamma$ MC shapes. Fig. 2, left, shows the fit in $M_{\text{miss}}^2$.

From the fit we obtain about 900 $\gamma\gamma \rightarrow \eta \rightarrow 3\pi^0$ events and a preliminary value for the cross section, measured with the $\eta \rightarrow \pi^0\pi^0\pi^0$ channel, $\sigma_{e^+e^- \rightarrow \eta\gamma}(1 \text{ GeV}) = 0.875 (18)_{\text{stat}} (35)_{\text{syst}}$, in agreement with that measured through the $\eta \rightarrow \pi^+\pi^-\pi^0$ channel. This is shown in Fig. 2 (right) among other experimental results [4] as a function of $\sqrt{s}$. We are finalizing the evaluation of the $\sigma_{\gamma\gamma \rightarrow \eta \rightarrow \pi^0\pi^0\pi^0}(1 \text{ GeV})$ cross section.

4. Search for $\gamma\gamma \rightarrow \pi^0\pi^0$

The main goal of this analysis is to investigate the low $\pi^0\pi^0$ invariant mass region, just above the production threshold, where a contribution from the $\sigma(600)$ scalar meson [7, 8, 9] as an intermediate state is expected.

The main backgrounds are annihilation processes with four or more prompt photons in the final state: $e^+e^- \rightarrow K_SK_L, \eta\gamma, \omega\pi^0, f_0(980)(\rightarrow \pi^0\pi^0)\gamma$. In addition, due to the possibility of cluster splitting, the $e^+e^- \rightarrow \gamma\gamma$ process is also considered as a source of background. Selected events have no tracks in the drift chamber and four prompt photons in the final state with polar angle $23^\circ < \theta < 157^\circ$ and energy $E_\gamma > 15 \text{ MeV}$. The photons are paired choosing the combination which minimizes the $\chi^2$-like variable of eq.(4) in the case of four photons coming from two neutral pions, $\chi^2_{4\gamma}$. Events with $\chi^2_{4\gamma} > 4$ are rejected: the effect of this selection is shown in Fig. 3 (left). To reject $e^+e^- \rightarrow K_SK_L$ events, where a large amount of non-prompt energy is released in the detector a tighter cut on the energy fraction carried by photons,
5. Conclusions

From an integrated luminosity of 240 pb$^{-1}$ of data collected at DAΦNE, operating at $\sqrt{s} \simeq 1\text{ GeV}$, the following preliminary results in $\gamma\gamma$ analyses are achieved:

- unambiguous signature of both $\gamma\gamma \to \eta$ and $\gamma\gamma \to \pi^0\pi^0$ events, without any $e^\pm$ tagger;
- $\gamma\gamma \to \eta$ events are observed through both $\eta \to \pi^+\pi^-\pi^0$ and $\eta \to \pi^0\pi^0\pi^0$ channels, with independent systematics;
- from the same data sample, an exploratory research shows a structure at small values of the $M_{4\gamma}$ spectrum, where the process $e^+e^- \to e^+e^-\pi^0\pi^0$ is expected.

As a by-product, we determined the cross section $\sigma_{e^+e^-\to\eta\gamma}$ at $\sqrt{s} = 1\text{ GeV}$, with accuracy better than the closer data points. These results are encouraging also in view of the forthcoming data taking campaign of the KLOE-2 project [10], when both low and high energy $e^\pm$ tagging devices will be available.

Acknowledgments

We thank the DAFNE team for their efforts in maintaining low background running conditions and their collaboration during all data-taking. We want to thank our technical staff: G. F. Fortugno and F. Sborzacchi for their dedication in ensuring efficient operation of the KLOE computing facilities; M. Anelli for his continuous attention to the gas system and detector safety; A. Balla, M. Gatta, G. Corradi and G. Papalino for electronics maintenance; M. Santoni, G. Paoluzzi and R. Rosellini for general detector support; C. Piscitelli for his help during major...
maintenance periods. This work was supported in part by EURODAPHNE, contract FMRX-CT98-0169; by the German Federal Ministry of Education and Research (BMBF) contract 06-KA-957; by the German Research Foundation (DFG), 'Emmy Noether Programme', contracts DE839/1-4; by the EU Integrated Infrastructure Initiative HadronPhysics Project under contract number RI3-CT-2004-506078; by the European Commission under the 7th Framework Programme through the 'Research Infrastructures' action of the 'Capacities' Programme, Call: FP7-INFRASTRUCTURES-2008-1, Grant Agreement N. 227431; by the Polish Ministry of Science and Higher Education through the Grant No. 0469/B/H03/2009/37.

References

[1] S. J. Brodsky, T. Kinoshita and H. Terazawa, Phys. Rev. D 4, 1532 (1971)
[2] F. Low, Phys. Rev. 120, 582 (1960)
[3] F. Nguyen, F. Piccinini and A. D. Polosa, Eur. Phys. J. C 47, 65 (2006)
[4] M. N. Achasov et al. [SND Collaboration], Phys. Rev. D 76, 077101 (2007)
[5] C. Amsler, T. Gutsche, S. Spanier and N. A. Törnqvist “Note on Scalar Mesons” in K. Nakamura et al. [Particle Data Group], J. Phys. G 37, 075021 (2010)
[6] H. Marsiske et al. [Crystal Ball Collaboration], Phys. Rev. D 41, 3324 (1990)
[7] E. M. Aitala et al. [E791 Collaboration], Phys. Rev. Lett. 86, 770 (2001)
[8] H. Muramatsu et al. [CLEO Collaboration], Phys. Rev. Lett. 89, 251802 (2002) [Erratum-ibid. 90, 059901 (2003)].
[9] M. Ablikim et al. [BES Collaboration], Phys. Lett. B 598, 149 (2004).
[10] G. Amelino-Camelia et al., Eur. Phys. J. C 68, 619 (2010)