Nonintrusive and structure preserving multiscale integration of stiff ODEs, SDEs and Hamiltonian systems with hidden slow dynamics via flow averaging

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April 14, 2010

Abstract

We introduce a new class of integrators for stiff ODEs as well as SDEs. Examples of subclasses of systems that we treat are ODEs and SDEs that are sums of two terms, one of which has large coefficients. These integrators are (i) Multiscale: they are based on flow averaging and so do not fully resolve the fast variables and have a computational cost determined by slow variables (ii) Versatile: the method is based on averaging the flows of the given dynamical system (which may have hidden slow and fast processes) instead of averaging the instantaneous drift of assumed separated slow and fast processes. This bypasses the need for identifying explicitly (or numerically) the slow or fast variables (iii) Nonintrusive: A pre-existing numerical scheme resolving the microscopic time scale can be used as a black box and easily turned into one of the integrators in this paper by turning the large coefficients on over a microscopic timescale and off during a mesoscopic timescale (iv) Convergent over two scales: strongly over slow processes and in the sense of measures over fast ones. We introduce the related notion of two-scale flow convergence and analyze the convergence of these integrators under the induced topology (v) Structure preserving: They inherit the structure preserving properties of the legacy integrators from which they are derived. Therefore, for stiff Hamiltonian systems (possibly on manifolds), they can be made to be symplectic, time-reversible, and symmetry preserving (symmetries are group actions that leave the system invariant) in all variables. They are explicit and applicable to arbitrary stiff potentials (that need not be quadratic). Their application to the Fermi-Pasta-Ulam problems shows accuracy and stability over four orders of magnitude of time scales. For stiff Langevin equations, they are symmetry preserving, time-reversible and Boltzmann-Gibbs reversible, quasi-symplectic on all variables and conformally symplectic with isotropic friction.

Contents

1 Overview of the integrator on ODEs
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Consider the following ODE on $\mathbb{R}^d$,

$$\dot{u}^\epsilon = G(u^\epsilon) + \frac{1}{\epsilon} F(u^\epsilon).$$  \hfill (1.1)

In Subsections 1.9, 2.1, 3.1, 3.5 and 4.1 we will consider more general ODEs, stiff deterministic Hamiltonian systems (2.1), SDEs ((3.1) and (3.15)) and Langevin equations ((4.1) and (4.2)); however for the sake of clarity, we will start the description of our method with (1.1).

**Condition 1.1.** Assume that there exists a diffeomorphism $\eta := (\eta^x, \eta^y)$, from $\mathbb{R}^d$ onto $\mathbb{R}^{d-p} \times \mathbb{R}^p$ (with uniformly bounded $C^1, C^2$ derivatives), separating slow and fast variables, i.e., such that (for all $\epsilon > 0$) the process $(x^\epsilon_t, y^\epsilon_t) = (\eta^x(u^\epsilon_t), \eta^y(u^\epsilon_t))$ satisfies an ODE system of the form

$$\begin{cases}
\dot{x}^\epsilon = g(x^\epsilon, y^\epsilon) & x^\epsilon_0 = x_0 \\
\dot{y}^\epsilon = \frac{1}{\epsilon} f(x^\epsilon, y^\epsilon) & y^\epsilon_0 = y_0.
\end{cases}$$  \hfill (1.2)

**Condition 1.2.** Assume that the fast variables in (1.2) are locally ergodic with respect to a family of measures $\mu$ drifted by slow variables. More precisely, we assume that there exists a family of probability measures $\mu(x, dy)$ on $\mathbb{R}^p$ indexed by $x \in \mathbb{R}^{d-p}$ and a positive function $T \mapsto E(T)$ such that $\lim_{T \to \infty} E(T) = 0$ and such that for all $x_0, y_0, T$ and $\phi$ uniformly bounded and Lipschitz, the solution to

$$\dot{Y}_t = f(x_0, Y_t) \quad Y_0 = y_0$$  \hfill (1.3)

satisfies

$$\left| \frac{1}{T} \int_0^T \phi(Y_s) ds - \int_{\mathbb{R}^p} \phi(y) \mu(x_0, dy) \right| \leq \chi(\|(x_0, y_0)\|) E(T)(\|\phi\|_{L^\infty} + \|\nabla \phi\|_{L^\infty})$$  \hfill (1.4)
where \( r \mapsto \chi(r) \) is bounded on compact sets.

Under conditions 1.1 and 1.2, it is known (we refer for instance to [95] or to Theorem 14, Section 3 of Chapter II of [104] or to [88]) that \( x^\epsilon \) converges towards \( x_t \) defined as the solution to the ODE

\[
\dot{x} = \int g(x, y) \mu(x, dy), \quad x|_{t=0} = x_0
\]  

where \( \mu(x, dy) \) is the ergodic measure associated with the solution to the ODE

\[
\dot{y} = f(x, y)
\]  

It follows that the slow behavior of solutions of (1.1) can be simulated over coarse time steps by first identifying the slow process \( x^\epsilon \) and then using numerical approximations of solutions of (1.2) to approximate \( x^\epsilon \). Two classes of integrators have been founded on this observation: The equation free method [64, 65] and the Heterogeneous Multiscale Method [36, 40, 35, 5]. One shared characteristic of the original form of those integrators is, after identification of the slow variables, to use a micro-solver to approximate the effective drift in (1.5) by averaging the instantaneous drift \( g \) with respect to numerical solutions of (1.6) over a time span larger than the mixing time of the solution to (1.6).

### 1.1 FLAVORS

In this paper, we propose a new method based on the averaging of the instantaneous flow of the ODE (1.1) with hidden slow and fast variables instead of the instantaneous drift of \( x^\epsilon \) in ODE (1.2) with separated slow and fast variables. We have called the resulting class of numerical integrators **Flow AVeraging integratORS (FLAVORS)**. Since FLAVORS are directly applied to (1.1), hidden slow variables do not need to be identified, either explicitly or numerically. Furthermore FLAVORS can be implemented using an arbitrary legacy integrator \( \Phi_{\frac{1}{h}} \) for (1.1) in which the parameter \( \frac{1}{h} \) can be controlled (figure 1). More precisely, assume that there exists a constant \( h_0 > 0 \) such that \( \Phi_{\alpha h} \) satisfies for all \( h \leq h_0 \min(\frac{1}{\alpha}, 1) \) and \( u \in \mathbb{R}^d \)

\[
|\Phi_{\alpha h}(u) - u - hG(u) - ahF(u)| \leq Ch^2(1 + \alpha)^2
\]  

then FLAVOR can be defined as the algorithm simulating the process

\[
\bar{u}_t = \left( \Phi_{\frac{1}{h} - \tau} \circ \Phi_{\frac{1}{h}} \right)^k(u_0) \quad \text{for} \quad k\delta \leq t < (k + 1)\delta
\]  

where \( \tau \) is a fine time step resolving the fast time scale \( (\tau \ll \epsilon) \) and \( \delta \) is a mesoscopic time step independent of the fast time scale satisfying \( \tau \ll \epsilon \ll \delta \ll 1 \) and

\[
\left( \frac{\tau}{\epsilon} \right)^2 \ll \delta \ll \frac{\tau}{\epsilon}
\]  

In our numerical experiments, we have used the “rule of thumb” \( \delta \sim \gamma \frac{\tau}{\epsilon} \) where \( \gamma \) is a small parameter (0.1 for instance).
A pre-existing numerical scheme resolving the microscopic time scale can be used as a black box and turned into a FLAVOR by simply turning on and off stiff parameters over a microscopic timescale $\tau$ (on) and a mesoscopic timescale $\delta$ (off). The bottom line of the approach is to (repeatedly) compose an accurate, short-time integration of the complete set of equations with an accurate, intermediate-time integration of the non-stiff part of the system. While the integration over short time intervals is accurate (in a strong sense), this is extended to intermediate time integration (in the sense of measures) using the interplay between the short time integration and the mesoscopic integration. The computational cost remains bounded independently from the stiff parameter $1/\epsilon$ because: (i) The whole system is only integrated over extremely short ($\tau \ll \epsilon$) time intervals during every intermediate ($\delta$) time intervals. (ii) The intermediate time step $\delta$ (that of the non-stiff part of the system) is limited not by the fast time scales ($\epsilon$) but by the slow ones ($\mathcal{O}(1)$).

By switching stiff parameters FLAVOR approximates the flow of (1.1) over a coarse time step $h$ (resolving the slow time scale) by the flow

$$
\Phi_h := \left( \Phi^{0}_{\frac{h}{\tau}} \circ \Phi^{1}_{\frac{1}{\epsilon}} \right)^M
$$

where $M$ is a positive integer corresponding to the number of “samples” used to average the flow ($\delta$ has to be identified with $\frac{h}{M}$). We refer to subsection 1.4 for the distinction between macro and meso-steps, for the rationale and mechanism behind FLAVORS and the limits (1.9).

Since FLAVORS are obtained by flow-composition, we will show in Section 2 and 4 that they inherit the structure preserving properties (for instance symplecticity and symmetries under a group action) of the legacy integrator for Hamiltonian systems and Langevin equations.

Under conditions (1.9) on $\tau$ and $\delta$, we show that (1.8) is strongly accurate with respect to (hidden) slow variables and weakly (in the sense of measures) accurate with respect to (hidden) fast variables. Motivated by this observation, we introduce the related notion of two-scale flow convergence in analogy with homogenization theory for elliptic PDEs [86, 3] and call it F-convergence for short. F-convergence is close in spirit to the Young measure approach to computing slowly advancing fast oscillations introduced in [10, 9].
1.2 Two-scale flow convergence

Let $(\xi^\epsilon_t)_{t \in \mathbb{R}^+}$ be a sequence of processes on $\mathbb{R}^d$ (functions from $\mathbb{R}^+$ to $\mathbb{R}^d$) indexed by $\epsilon > 0$. Let $(X_t)_{t \in \mathbb{R}^+}$ be a process on $\mathbb{R}^{d-p}$ ($p \geq 0$). Let $x \mapsto \nu(x, dz)$ be a function from $\mathbb{R}^{d-p}$ into the space of probability measures on $\mathbb{R}^d$.

**Definition 1.1.** We say that the process $\xi^\epsilon_t F$-converges to $\nu(X_t, dz)$ as $\epsilon \downarrow 0$ and write $\xi^\epsilon_t \overset{F}{\longrightarrow} \nu(X_t, dz)$ if and only if for all functions $\varphi$ bounded and uniformly Lipshitz-continuous on $\mathbb{R}^d$, and for all $t > 0$,

$$
\lim_{h \to 0} \lim_{\epsilon \to 0} \frac{1}{h} \int_t^{t+h} \varphi(\xi^\epsilon_s) \, ds = \int_{\mathbb{R}^d} \varphi(z) \nu(X_t, dz)
$$

(1.11)

1.3 Asymptotic convergence result

Our convergence theorem requires that $u^\epsilon_t$ and $\bar{u}_t$ do not blow up as $\epsilon \downarrow 0$; more precisely, we will assume that the following conditions are satisfied.

**Condition 1.3.** Assume that:

1. $F$ and $G$ are Lipschitz continuous.
2. For all $u_0$, $T > 0$, the trajectories $(u^\epsilon_t)_{0 \leq t \leq T}$ are uniformly bounded in $\epsilon$.
3. For all $u_0$, $T > 0$, the trajectories $(\bar{u}_t^\epsilon)_{0 \leq t \leq T}$ are uniformly bounded in $\epsilon$, $0 < \delta \leq h_0$, $\tau \leq \min(\tau_0 \epsilon, \delta)$.

For $\pi$, an arbitrary measure on $\mathbb{R}^d$, we define $\eta^{-1} * \pi$ to be the push forward of the measure $\pi$ by $\eta^{-1}$.

**Theorem 1.1.** Let $u^\epsilon_t$ be the solution to (1.1) and $\bar{u}_t$ be defined by (1.8). Assume that equation (1.7) and conditions 1.1, 1.2 and 1.3 are satisfied, then

- $u^\epsilon_t$ F-converges to $\eta^{-1} * (\delta_X \otimes \mu(X_t, dy))$ as $\epsilon \downarrow 0$ where $X_t$ is the solution to

$$
\dot{X}_t = \int g(X_t, y) \mu(X_t, dy) \quad X_0 = x_0.
$$

(1.12)

- $\bar{u}_t$ F-converges to $\eta^{-1} * (\delta_X \otimes \mu(X_t, dy))$ for $\epsilon \leq \delta / (-C \ln \delta)$, $\frac{\tau}{\epsilon} \downarrow 0$, $\frac{\tau}{\epsilon} \delta \downarrow 0$ and $(\frac{\tau}{\epsilon})^2 \frac{1}{\delta} \downarrow 0$.

**Remark 1.1.** The $F$-convergence of $u^\epsilon_t$ to $\eta^{-1} * (\delta_X \otimes \mu(X_t, dy))$ can be restated as

$$
\lim_{h \to 0} \lim_{\epsilon \to 0} \frac{1}{h} \int_t^{t+h} \varphi(u^\epsilon_s) \, ds = \int_{\mathbb{R}^d} \varphi(\eta^{-1}(X_t, y)) \mu(X_t, dy)
$$

(1.13)

for all functions $\varphi$ bounded and uniformly Lipshitz-continuous on $\mathbb{R}^d$, and for all $t > 0$. 6
Remark 1.2. Observe that $g$ comes from (1.5). It is not explicitly known and does not need to be explicitly known for the implementation of the proposed method.

Remark 1.3. The limits on $\epsilon, \tau$ and $\delta$ are in essence stating that FLAVOR is accurate provided that $\tau \ll \epsilon$ ($\tau$ resolves the stiffness of (1.1)) and equation (1.9) is satisfied.

Remark 1.4. Throughout this paper, $C$ will refer to an appropriately large enough constant independent from $\epsilon, \delta, \tau$. To simplify the presentation of our results, we use the same letter $C$ for expressions such as $2CeC$ instead of writing it as a new constant $C_1$ independent from $\epsilon, \delta, \tau$.

1.4 Rationale and mechanism behind FLAVORS

We will now explain the rationale and mechanism behind FLAVORS. We refer to Subsection [7.1] of the appendix for the detailed proof of Theorem [1.1]. Let us start by considering the case where $\eta$ is the identity diffeomorphism. Let $\varphi^1_{1/\epsilon}$ be the flow of (1.2).

Observe that $\varphi^0_0$ (obtained from $\varphi^1_{1/\epsilon}$ by setting the parameter $1/\epsilon$ to zero) is the flow of (1.2) with $y^\epsilon$ frozen, i.e.,

$$
\varphi^0_0(x, y) = (\hat{x}_t, y) \quad \text{where } \hat{x}_t \text{ solves } \frac{d\hat{x}}{dt} = g(\hat{x}, y), \quad \hat{x}_0 = x.
$$

The main effect of FLAVORS is to average the flow of (1.2) with respect to fast degrees of freedom via splitting and re-synchronization. By splitting, we refer to the substitution of the flow $\varphi^1_{1/\epsilon}$ by composition of $\varphi^0_0$ and $\varphi^1_{1/\epsilon}$, and by re-synchronization we refer to the distinct time-steps $\delta$ and $\tau$ whose effects are to advance the internal clock of fast variables by $\tau$ every step of length $\delta$. By averaging, we refer to the fact that FLAVORS approximates the flow $\varphi^1_{1/\epsilon}$ by the flow

$$
\varphi_h := (\varphi^0_0 \circ \varphi^1_{1/\epsilon})^M
$$

where $h$ is a coarse time step resolving the slow time scale associated with $x^\epsilon$, $M$ is a positive integer corresponding to the number of samples used to average the flow (where $\delta$ is identified with $h/M$) and $\tau$ is a fine time step resolving the fast time scale, of the order of $\epsilon$, and associated with $y^\epsilon$. In general, analytical formulae are not available for $\varphi^0_0$ and $\varphi^1_{1/\epsilon}$ and numerical approximations are used instead.

Observe that when FLAVORS are applied to systems with explicitly separated slow and fast processes, they lead to integrators that are locally in the neighborhood of those obtained with HMM (or equation free) methods with a reinitialization of the fast variables at macrotime $n$ by their final value at macrotime step $n - 1$ and with only one microstep per macrostep [37, 39].

We will now consider the situation where $\eta$ is not the identity diffeomorphism and
give the rationale behind the limits (1.9).

\[
\begin{array}{c}
\eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \\
\eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \xrightarrow{\Phi^0_{\delta-\tau}} \eta^{-1} \\
\end{array}
\]

As illustrated in the above diagram, since \((\bar{x}_t, \bar{y}_t) = \eta(\bar{u}_t)\), simulating \(\bar{u}_{n\delta}\) defined in (1.8) is equivalent to simulating the discrete process

\[
(\bar{x}_{n\delta}, \bar{y}_{n\delta}) := \left( \Psi^1_{\delta-\tau} \circ \Psi^0_{\tau} \right)^n(x_0, y_0) \quad (1.16)
\]

where

\[
\Psi^0_h := \eta \circ \Phi^0_h \circ \eta^{-1} \quad (1.17)
\]

Observe that the accuracy (in the topology induced by F-convergence) of \(\bar{u}_t\) with respect to \(u^t\), solution of (1.1), is equivalent to that of \((\bar{x}_t, \bar{y}_t)\) with respect to \((x^t, y^t)\) defined by (1.2). Now, for the clarity of the presentation, assume that \(\Phi^0_h(u) = u + hG(u) + \alpha hF(u)\) (1.18)

Using Taylor’s theorem and (1.18), we obtain that

\[
\Psi^1_h(x, y) = (x, y) + h(g(x, y), 0) + \alpha h(0, f(x, y)) + \int_0^1 v^T \text{Hess } \eta(u) \tau v(1-t)^2 \, dt \quad (1.19)
\]

with

\[
u := \eta^{-1}(x, y) \quad \text{and} \quad v := h(G + \alpha F) \circ \eta^{-1}(x, y) \quad (1.20)
\]

It follows from equations (1.19) and (1.20) that \(\Psi^1_h\) is a first order accurate integrator approximating the flow of (1.2) and \(\Psi^0_h\) is a first order accurate integrator approximating the flow of (1.14). Let \(h\) be a coarse time step and \(\delta\) a mesostep. Since \(\bar{x}\) remains nearly constant over the coarse time step, the switching (on and off) of the stiff parameter \(\frac{1}{\epsilon}\) averages the drift \(g\) of \(\bar{x}\) with respect to the trajectory of \(\bar{y}\) over \(h\). Since the coarse step \(h\) is composed of \(\frac{h}{\delta}\) mesosteps, the internal clock of the fast process is advanced by \(\frac{h}{\delta} \times \frac{\tau}{\epsilon}\). Since \(h\) is of the order of one, the trajectory of \(\bar{y}\) is mixing with respect to the local ergodic measure \(\mu\) provided that \(\frac{\tau}{\delta} \gg 1\), i.e.

\[
\delta \ll \frac{\tau}{\epsilon} \quad (1.21)
\]

Equation (1.21) corresponds to the right hand side of equation (1.9). If \(\eta\) is a non-linear diffeomorphism (with non-zero Hessian), it also follows from equations (1.19) and (1.20) that each invocation of the integrator \(\Psi^1_{\tau/\epsilon}\) occasions an error (on the accuracy of the slow process) proportional to \(\left(\frac{\tau}{\epsilon}\right)^2\). Since during the coarse time step \(h\), \(\Psi^1_{\tau/\epsilon}\) is solicited
$h$-times, it follows that the error accumulation during $h$ is $\frac{h}{\delta} \times (\frac{\tau}{\epsilon})^2$. Hence, the accuracy of the integrator requires that $\frac{h}{\delta} \times (\frac{\tau}{\epsilon})^2 \ll 1$, i.e.

$$\left(\frac{\tau}{\epsilon}\right)^2 \ll \delta$$

(1.22)

Equation (1.22) corresponds to the left hand side of equation (1.9).

Observe that if $\eta$ is linear, its Hessian is null and the remainder in the right hand side of (1.19) is zero. It follows that if $\eta$ is linear, the error accumulation due to fine time steps on slow variables is zero and condition (1.21) is sufficient for the accuracy of the integrator.

It has been observed in [38] and in Section 5 of [111] that slow variables do not need to be identified with HMM/averaging type integrators if the relation between original and slow variables is linear or a permutation and if

$$M \Delta t \ll \frac{\tau}{\epsilon}$$

(1.23)

where $M$ the number of fine-step iterations used by HMM to compute the average the drift of slow variables and $\Delta t$ is the coarse time step (in HMM) along the direction of the averaged drift. The analysis of FLAVORS associated with equation (1.19) reaches a similar conclusion if $\eta$ is linear in the sense that the error caused by the Hessian of $\eta$ in (1.19) is zero and in the (sufficient) condition (1.21) is analogous to (1.23) for $M = 1$. It is also stated on Page 2 of [38] that “there are counterexamples showing that algorithms of the same spirit do not work for deterministic ODEs with separated time scales if the slow variables are not explicitly identified and made use of. But in the present context, the slow variables are linear functions of the original variables, and this is the reason why the seamless algorithm works.” Here, the analysis of FLAVORS associated with equation (1.19) shows an algorithm based on an averaging principle would indeed, in general, not work if $\eta$ is nonlinear (and (1.22) not satisfied) due to the error accumulation (on slow variables) associated with the Hessian of $\eta$. However, the above analysis also shows that if condition (1.22) is satisfied, then, although $\eta$ may be nonlinear, flow averaging integrators will always work without identifying slow variables.

1.5 Non asymptotic convergence result

**Theorem 1.2.** Under assumptions and notations of theorem 1.1, there exists $C > 0$ such that for $\delta < h_0$, $\tau < \tau_0 \epsilon$ and $t > 0$,

$$|x_t^\epsilon - \eta^T(\bar{u}_t)| \leq C e^{Ct} \chi_1(u_0, \epsilon, \delta, \tau)$$

(1.24)

and

$$\left| \frac{1}{T} \int_t^{t+T} \varphi(\bar{u}_s) \, ds - \int_{\mathbb{R}^p} \varphi(\eta^{-1}(X_t, y)) \mu(X_t, dy) \right|$$

$$\leq \chi_2(u_0, \epsilon, \delta, \tau, T, t)(\|\varphi\|_{L^\infty} + \|\nabla \varphi\|_{L^\infty})$$

(1.25)

where $\chi_1$ and $\chi_2$ are functions converging towards zero as $\epsilon \leq \delta/(C \ln \frac{1}{\delta})$, $\frac{\tau}{\epsilon} \downarrow 0$, $\frac{\tau}{\epsilon} \delta \downarrow 0$ and $(\frac{\tau}{\epsilon})^2 \downarrow 0$ (and $T \downarrow 0$ for $\chi_2$).
Remark 1.5. For \( \epsilon \leq \delta / ( -C \ln \delta) \) and \( \delta \frac{\tau}{\epsilon} + \frac{\tau}{\epsilon} \leq 1 \), the following holds
\[
\chi_1(u_0, \epsilon, \delta, \tau) \leq \sqrt{\delta} + \left( \frac{\tau}{\epsilon} \right)^{1/2} + E \left( \frac{1}{C} \ln \frac{1}{\delta} \right) + \left( \frac{\delta \epsilon}{\tau} \right)^{1/2} + E \left( \frac{1}{C} \ln \left( \frac{\delta \epsilon + \tau}{\epsilon} \right)^{-1} \right)
\]
(1.26)
and \( \chi_2 \) satisfies a similar inequality.

Remark 1.6. Choosing \( \tau \sim \gamma \epsilon \) and \( \delta \sim \gamma \frac{\tau}{\epsilon} \), where \( \gamma \) is a small constant independent from \( \epsilon \), Theorem 1.2 shows that the approximation error of FLAVOR is bounded by a function of \( \gamma \) converging towards zero as \( \gamma \downarrow 0 \). If follows that the speed up is of the order of \( \frac{\delta}{\epsilon} \sim \frac{\tau}{\epsilon} \), i.e., scales like \( \frac{1}{\epsilon} \) at fixed accuracy. In order to be able to compare FLAVOR with integrators resolving all the fine time steps we have limited the speed up in the numerical experiments to 200\( \times \) (but the latter can be arbitrary large as \( \epsilon \downarrow 0 \)). For sufficiently small \( \epsilon \), we observe that FLAVORS with microstep \( \tau \) and mesostep \( \delta \) overperform their associated legacy integrator with the same microstep \( \tau \) over large simulation times (we refer to Section 6.3 on the Fermi-Pasta-Ulam problem). This phenomenon is caused by an error accumulation at each tick (microstep) of the clock of fast variables. Since FLAVORS (indirectly, i.e., without identifying fast variables) slow down the speed of this clock from \( \frac{1}{\epsilon} \) to a value \( \frac{\tau}{\epsilon} \sim \frac{1}{\epsilon} \) independent from \( \epsilon \) this error does not blow up as \( \epsilon \downarrow 0 \) (as opposed to an integrator resolving all fine time steps). Because of this reason, if this error accumulation on fast variables is exponential, then the speed up at fixed accuracy does not scale like \( \frac{1}{\epsilon} \), but like \( e^{-T \epsilon} \) where \( T \) is the total simulation time. A consequence of this phenomenon can be seen in Figure 10 (associated with the FPU problem) where Velocity Verlet fails to capture the \( O(\epsilon^{-1}) \) dynamics with a time step \( h = 10^{-5} \) whereas FLAVORS remain accurate with \( \tau = 10^{-4} \) and \( \delta = 2 \cdot 10^{-3} \).

Remark 1.7. The reader should not be surprised by the presence of the exponential factor \( e^{Ct} \) in (1.24). It is known that global errors for numerical approximations of ODEs grow, in general, exponentially with time (see for instance [55]). These bounds are, however, already tight; consider, for instance, how error propagates in a generic numerical scheme applied to the special system of \( \dot{x} = x \). It is possible to show that the increase of global errors is linear in time only for a restricted class of ODEs (using techniques from Lyapunov’s theory of stability [115]). Notice that the constant \( C \) in the exponential of our bound does not scale with \( \epsilon^{-1} \), and therefore the bound is uniform and rather tight.

Remark 1.8. We refer to [40] for higher order averaging based methods. In particular, [40] shows how, after identification of slow variables, balancing the different error contributions yields an explicit stable integration method having the order of the macro scheme.

1.6 Natural FLAVORS

Although convenient, it is not necessary to use legacy integrators to obtain FLAVORS. More precisely, theorems 1.1 and 1.2 remain valid if FLAVORS are defined to be algo-
rithms simulating the discrete process

\[ \bar{u}_t := (\theta_{k+1}^{\epsilon} \circ \theta_{k+1}^{G})^k (u_0) \quad \text{for} \quad k\delta \leq t < (k+1)\delta \tag{1.27} \]

where \( \theta_{k+1}^{\epsilon} \) and \( \theta_{k+1}^{G} \) are two mappings from \( \mathbb{R}^d \) onto \( \mathbb{R}^d \) (the former approximating the flow of the whole system (1.1) for time \( \tau \), and the latter approximating the flow of \( \dot{v} = G(v) \) for time \( \delta - \tau \)), satisfying the following conditions.

**Condition 1.4.** Assume that:

1. There exists \( h_0, C > 0 \) such that for \( h \leq h_0 \) and any \( u \in \mathbb{R}^d \),

\[ |\theta_{h}^{G}(u) - u - hG(u)| \leq Ch^2 \tag{1.28} \]

2. There exists \( \tau_0, C > 0 \), such that for \( \tau \epsilon \leq \tau_0 \) and any \( u \in \mathbb{R}^d \),

\[ |\theta_{\tau}^{\epsilon}(u) - u - \tau G(u) - \tau \epsilon F(u)| \leq C (\tau \epsilon)^2 \tag{1.29} \]

3. For all \( u_0, T > 0 \), the discrete trajectories \( (\theta_{k+1}^{G} \circ \theta_{k}^{\epsilon})^k (u_0) \) are uniformly bounded in \( \epsilon \), \( 0 < \delta \leq h_0 \), \( \tau \leq \min(\tau_0, \delta) \).

Observe that (1.8) is a particular case of (1.27) in which \( \theta^{\epsilon} = \Phi^{1/\epsilon} \) and the mapping \( \theta^{G} \) is obtained from the legacy integrator \( \Phi^{\alpha} \) by setting \( \alpha \) to zero.

1.7 Related work

Dynamical systems with multiple time scales pose a major problem in simulations because the small time steps required for stable integration of the fast motions lead to large numbers of time steps required for the observation of slow degrees of freedom [109, 54]. Traditionally, stiff dynamical systems have been separated into two classes with distinct integrators: stiff systems with fast transients and stiff systems with rapid oscillations [6, 35, 98]. The former has been solved using implicit schemes [16, 34, 51, 56], Chebyshev methods [70, 1] or the projective integrator approach [48]. These latter have been solved using filtering techniques [45, 66, 100] or Poincaré map techniques [47, 90]. We also refer to methods based on highly oscillatory quadrature [32, 60, 59], an area that has undergone significant developments in the last few years [61]. It has been observed that at the present time, there exists no unified strategy for dealing with both classes of problems [35]. When slow variables can be identified, effective equations can be obtained by averaging the instantaneous drift driving those slow variables [104]. Two classes of numerical methods have been built on this observation: The equation-free method [64, 65] and the Heterogeneous Multiscale Method [36, 40, 35, 5]. Observe that FLAVORS apply in a unified way to both stiff systems with fast transients and stiff systems with rapid oscillations, with or without noise, with a mesoscopic integration time step chosen independently from the stiffness.
1.8 Limitations of the method

The proof of the accuracy of the method (theorems 1.1 and 1.2) is based on an averaging principle; hence, if \( \epsilon \) is not small (the stiffness of the ODE is weak), although the method may be stable, there is no guarantee of accuracy. More precisely, the global error of the method is an increasing function of \( \epsilon, \delta, \frac{\delta}{\epsilon}, (\frac{\delta}{\epsilon})^2 \). Writing \( \gamma := \frac{\delta}{\epsilon} \) the accuracy method requires \( \gamma^2 \ll \delta \ll \gamma \). Choosing \( \delta = \gamma^3 \), the condition \( \epsilon \ll \delta \ll 1 \) (related to computational gain) requires \( \epsilon^{\frac{3}{2}} \ll \gamma \ll 1 \) which can be satisfied only if \( \epsilon \) is small.

The other limitation of the method lies in the fact that a stiff parameter \( \frac{1}{\epsilon} \) needs to be clearly identified. In many examples of interest (Navier-Stokes equations, Maxwell’s equations,...), stiffness is a result of nonlinearity, initial conditions or boundary conditions and not of the existence of a large parameter \( \frac{1}{\epsilon} \). Molecular dynamics can also create widely separated time-scales from non-linear effects; we refer, for instance, to [116] and references therein.

1.9 Generic stiff ODEs

FLAVORS have a natural generalization to systems of the form

\[
\dot{u}^{\alpha,\epsilon} = F(u^{\alpha,\epsilon}, \alpha, \epsilon) \tag{1.30}
\]

where \( u \mapsto F(u, \alpha, \epsilon) \) is Lipshitz continuous.

**Condition 1.5.** Assume that:

1. \( \epsilon \mapsto F(u, \alpha, \epsilon) \) is uniformly continuous in the neighborhood of 0.

2. There exists a diffeomorphism \( \eta := (\eta^x, \eta^y) \), from \( \mathbb{R}^d \) onto \( \mathbb{R}^{d-p} \times \mathbb{R}^p \), independent from \( \epsilon, \alpha \), with uniformly bounded \( C^1, C^2 \) derivatives, such that the process \( (x^\alpha_t, y^\alpha_t) = (\eta^x(u^\alpha_t,0), \eta^y(u^\alpha_t,0)) \) satisfies, for all \( \alpha \geq 1 \), the ODE

\[
\dot{x}^\alpha = g(x^\alpha, y^\alpha) \quad x^\alpha_0 = x_0, \tag{1.31}
\]

where \( g(x,y) \) is Lipshitz continuous in \( x \) and \( y \) on bounded sets.

3. There exists a family of probability measures \( \mu(x, dy) \) on \( \mathbb{R}^p \) such that for all \( x_0, y_0, T \) \( (x_0, y_0) := \eta(u_0) \) and \( \varphi \) uniformly bounded and Lipschitz

\[
\left| \frac{1}{T} \int_0^T \varphi(y^\alpha_t) \, ds - \int_{\mathbb{R}^p} \varphi(y) \mu(dy) \right| \leq \chi(\|x_0, y_0\|)(E_1(T) + E_2(T \alpha^r)) \|\nabla \varphi\|_{L^\infty} \tag{1.32}
\]

where \( r \mapsto \chi(r) \) is bounded on compact sets and \( E_2(r) \to 0 \) as \( r \to \infty \) and \( E_1(r) \to 0 \) as \( r \to 0 \).

4. For all \( u_0, T > 0 \), the trajectories \( (u^\alpha_t)_{0 \leq t \leq T} \) are uniformly bounded in \( \alpha \geq 1 \).
Remark 1.9. Observe that slow variables are not kept frozen in equation (1.32). The error on local invariant measures induced by the (slow) drift of \(x^\alpha\) is controlled by \(E_2\). More precisely, the convergence of the right hand side of (1.32) towards zero requires the convergence of \(T\) towards zero and (at the same time) the divergence of \(T\alpha\) towards infinity.

Assume that we are given a mapping \(\Phi_{h,\alpha,\epsilon}\) from \(\mathbb{R}^d\) onto \(\mathbb{R}^d\) approximating the flow of (1.30). If the parameter \(\alpha\) can be controlled then \(\Phi_{h,\alpha,\epsilon}\) can be used as a black box for accelerating the computation of solutions of (1.30).

Condition 1.6. Assume that:

1. There exists a constant \(h_0 > 0\) such that \(\Phi_{h,\alpha,\epsilon}\) satisfies for all \(h \leq h_0\min(\frac{1}{\alpha^2}, 1)\),
\[
0 < 1 < \epsilon < 1 < \alpha
\]
\[
|\Phi_{h,\alpha,\epsilon}(u) - u - hF(u, \alpha, \epsilon)| \leq C(u)h^2(1 + \alpha^{2\epsilon}) \tag{1.33}
\]

where \(C(u)\) is bounded on compact sets.

2. For all \(u_0\), \(T > 0\), the discrete trajectories \(\left(\Phi_{\delta,\tau} \circ \Phi_{\tau}^{1/\epsilon}\right)(u_0)\) are uniformly bounded in \(0 < \epsilon \leq 1\), \(0 < \delta \leq h_0\), \(\tau \leq \min(h_0\epsilon^2, \delta)\).

FLAVOR can be defined as the algorithm given by the process
\[
\bar{u}_t = \left(\Phi_{\delta,\tau}^{1/\epsilon} \circ \Phi_{\tau}^{1/\epsilon}\right)(u_0) \text{ for } k\delta \leq t < (k + 1)\delta \tag{1.34}
\]

The theorem below shows the accuracy of FLAVORS for \(\delta \ll h_0\), \(\tau \ll \epsilon^\nu\) and \((\frac{\tau}{\epsilon^\nu})^2 \ll \delta \ll \frac{\tau}{\epsilon^\nu}\).

Theorem 1.3. Let \(u_{t,\epsilon}^{1/\epsilon}\) be the solution to (1.30) with \(\alpha = 1/\epsilon\) and \(\bar{u}_t\) be defined by (1.34). Assume that Conditions 1.5 and 1.6 are satisfied then:

- \(u_{t,\epsilon}^{1/\epsilon}\) F-converges towards \(\eta^{-1} * (\delta X_t \otimes \mu(X_t, dy))\) as \(\epsilon \downarrow 0\) where \(X_t\) is the solution to
\[
\dot{X}_t = \int_{\mathbb{R}^p} g(X_t, y) \mu(X_t, dy) \quad X_0 = x_0, \tag{1.35}
\]

- As \(\epsilon \downarrow 0\), \(\tau \epsilon^{-\nu} \downarrow 0\), \(\delta \epsilon^{-\nu} \downarrow 0\), \(\frac{\tau^2}{\epsilon^2 \Delta} \downarrow 0\), \(\bar{u}_t\) F-converges towards \(\eta^{-1} * (\delta X_t \otimes \mu(X_t, dy))\) as \(\epsilon \downarrow 0\) where \(X_t\) is the solution of (1.35).

Proof. The proof of Theorem 1.3 is similar to that of Theorem 1.1 and 3.1. Only the idea of the proof will be given here. The condition \(\epsilon \ll 1\) is needed for the approximation of \(u_{\alpha,\epsilon}\) by \(u_{\alpha,0}\) and for the F-convergence of \(u_{\epsilon,\epsilon}^{1/\epsilon}\). Since \(y_t^\alpha = \eta(h(u_{t,\epsilon}^{1/\epsilon}))\) the condition \(\tau \ll \epsilon^\nu\) is used along with equation (1.33) for the accuracy of \(\Phi_{\tau}^{1/\epsilon}\) in (locally) approximating \(y_t^\alpha\). The condition \((\frac{\tau}{\epsilon^\nu})^2 \ll \delta\) is required in order to control the error accumulated by \(m\) iterations of \(\Phi_{\tau}^{1/\epsilon}\). 

\[
\square
\]
2 Deterministic mechanical systems: Hamiltonian equations

Since averaging with FLAVORS is obtained by flow composition, FLAVORS have an inherent extension to multiscale structure preserving integrators for stiff Hamiltonian systems, i.e. ODEs of the form

\[ \dot{p} = -\partial_q H(p,q) \quad \dot{q} = \partial_p H(p,q) \] (2.1)

where the Hamiltonian

\[ H(q,p) := \frac{1}{2} p^T M^{-1} p + V(q) + \frac{1}{\epsilon} U(q) \] (2.2)

represents the total energy of a mechanical system with Euclidean phase space \( \mathbb{R}^d \times \mathbb{R}^d \) or a cotangent bundle \( T^*M \) of a configuration manifold \( M \).

Structure preserving numerical methods for Hamiltonian systems have been developed in the framework of geometric numerical integration \[54, 72\] and variational integrators \[78, 75\]. The subject of geometric numerical integration deals with numerical integrators that preserve geometric properties of the flow of a differential equation, and it explains how structure preservation leads to an improved long-time behavior \[53\]. Variational integration theory derives integrators for mechanical systems from discrete variational principles and are characterized by a discrete Noether theorem. These methods have excellent energy behavior over long integration runs because they are symplectic, i.e., by backward error analysis, they simulate a nearby mechanical system instead of nearby differential equations. Furthermore, statistical properties of the dynamics such as Poincaré sections are well preserved even with large time steps \[16\]. Preservation of structures is especially important for long time simulations. Consider integrations of a harmonic oscillator, for example: no matter how small a time step is used, the amplitude given by Forward Euler / Backward Euler will increase / decrease unboundedly, whereas the amplitude given by Variational Euler (also known as symplectic Euler) will be oscillatory with a variance controlled by the step length.

These long term behaviors of structure-preserving numerical integrators motivated their extension to multiscale or stiff Hamiltonian systems. We refer to \[31\] for a recent review on numerical integrators for highly oscillatory Hamiltonian systems. Symplectic integrators are natural for the integration of Hamiltonian systems since they reproduce at the discrete level an important geometric property of the exact flow \[20\]. For symplectic integrators primarily for (but not limited to) stiff quadratic potentials, we refer to the Impulse Method, the Mollified Impulse Method, and their variations \[51, 110, 44, 96\], which require an explicit form of the flow map of stiff process. In the context of variational integrators, by defining a discrete Lagrangian with an explicit trapezoidal approximation of the soft potential and a midpoint approximation for the fast potential, a symplectic (IMEX—IMplicit–EXplicit) scheme for stiff Hamiltonian systems has been proposed in \[105\]. The resulting scheme is explicit for quadratic potentials and implicit for non quadratic stiff potentials. We also refer to Le Bris and Legoll’s (Hamilton-Jacobi
derived) homogenization method [20]. Asynchronous Variational Integrators [74] provide a way to derive conservative symplectic integrators for PDEs where the solution advances non-uniformly in time; however, stiff potentials require a fine time step discretization over the whole time evolution. In addition, multiple time-step methods [106] evaluate forces to different extents of accuracies by approximating less important forces via Taylor expansions, but it has issues on long time behavior, stability and accuracy, as described in Section 5 of [73]. Fixman froze the fastest bond oscillations in polymers to remove stiffness by adding a log term resemblance of entropy-based free energy to compensate [43]. This approach is successful in studying statistics of the system, but does not always reconstruct the correct dynamics [91, 89, 14].

Several approaches to the homogenization of Hamiltonian systems (in analogy with classical homogenization [11, 62]) have been proposed. We refer to $\mathcal{M}$-convergence introduced in [101, 15], to the two-scale expansion of solutions of the Hamilton-Jacobi form of Newton’s equations with stiff quadratic potentials [20] and to PDE methods in weak KAM theory [41]. We also refer to [26], [58] and [95].

Obtaining explicit symplectic integrators for Hamiltonian systems with non-quadratic stiff potentials is known to be an important and nontrivial problem. By using Verlet/leapfrog macro-solvers, methods that are symplectic on slow variables (when those variables can be identified) have been proposed in the framework of HMM (the Heterogeneous Multiscale Method) in [102, 24]. A “reversible averaging” method has been proposed in [71] for mechanical systems with separated fast and slow variables. More recently, a reversible multiscale integration method for mechanical systems was proposed in [6] in the context of HMM. By tracking slow variables, [6] enforces reversibility in all variables as an optimization constraint at each coarse step when minimizing the distance between the effective drift obtained from the micro-solver (in the context of HMM) and the drift of the macro-solver. We are also refer to [99] for HMM symmetric methods for mechanical systems with a stiff potentials of the form $\frac{1}{\epsilon} \sum_{j=1}^{\nu} g_j(q)^2$.

2.1 FLAVORS for mechanical systems on manifolds

Assume that we are given a first order accurate legacy integrator for (2.1) in which the parameter $1/\epsilon$ can be controlled, i.e. a mapping $\Phi^0_h$ acting on the phase space such that for $h \leq h_0 \min(1, \alpha^{-\frac{1}{2}})$

$$\left| \Phi^0_h(q,p) - (q,p) - h(M^{-1}p, -V(q) - \alpha U(q)) \right| \leq Ch^2(1 + \alpha) \quad (2.3)$$

Write $\Theta_\delta$, the FLAVOR discrete mapping approximating solutions of (2.1) over time steps $\delta \gg \epsilon$, i.e.

$$(q_{(n+1)\delta}, p_{(n+1)\delta}) := \Theta_\delta(q_{n\delta}, p_{n\delta}). \quad (2.4)$$

FLAVOR can then be defined by

$$\Theta_\delta := \Phi^0_{\delta - \tau} \circ \Phi^{\frac{1}{\tau}}_\tau \quad (2.5)$$

Theorem 1.3 establishes the accuracy of this integrator under Conditions 1.5 and 1.6 provided that $\tau \ll \sqrt{\epsilon} \ll \delta$ and $\frac{\tau^2}{\epsilon} \ll \delta \ll \frac{\tau}{\sqrt{\epsilon}}$. 

15
2.1.1 Structure preserving properties of FLAVORS

We will now show that FLAVORS inherit the structure preserving properties of their legacy integrators.

**Theorem 2.1.** If for all \( h, \epsilon > 0 \) \( \Phi^\epsilon_h \) is symmetric under a group action, then \( \Theta^\delta \) is symmetric under the same group action.

**Theorem 2.2.** If \( \Phi^\alpha_h \) is symplectic on the co-tangent bundle \( T^*M \) of a configuration manifold \( M \), then \( \Theta^\delta \) defined by (2.5) is symplectic on the co-tangent bundle \( T^*M \).

Theorem 2.1 and Theorem 2.2 can be resolved by noting that “the overall method is symplectic - as a composition of symplectic transformations, and it is symmetric - as a symmetric composition of symmetric steps” (see Chapter XIII.1.3 of [54]).

Write

\[ \Phi^\epsilon_h := (\Phi_{-h})^{-1} \]  

(2.6)

Let us recall the following definition corresponding to definition 1.4 of the Chapter V of [54]

**Definition 2.1.** A numerical one-step method \( \Phi_h \) is called time-reversible if it satisfies \( \Phi_h^* = \Phi_h \).

The following theorem, whose proof is straightforward, shows how to derive a “symplectic and symmetric and time-reversible” FLAVOR from a symplectic legacy integrator and its adjoint. Since this derivation applies to manifolds, it also leads to structure-preserving FLAVORS for constrained mechanical systems.

**Theorem 2.3.** If \( \Phi^\alpha_h \) is symplectic on the co-tangent bundle \( T^*M \) of a configuration manifold \( M \), then

\[ \Theta^\delta := \Phi_{1/2} \circ \Phi_{-1/2} \circ \Phi_{\tau_2} \circ \Phi_{0} \circ \Phi_{\tau_2} \circ \Phi_{0} \circ \Phi_{\tau_2} \]  

(2.7)

is symplectic and time-reversible on the co-tangent bundle \( T^*M \).

**Remark 2.1.** Observe that (except for the first and last steps) iterating \( \Theta^\delta \) defined by (2.7) is equivalent to iterating

\[ \Theta^\delta := \Phi_{1/2} \circ \Phi_{-1/2} \circ \Phi_{\tau_2} \circ \Phi_{0} \circ \Phi_{\tau_2} \circ \Phi_{1/2} \]  

(2.8)

It follows that a symplectic, symmetric and reversible FLAVOR can be obtained in a nonintrusive way from a Störmer/Verlet integrator for (2.1) [53, 55, 114].

2.1.2 An example of a symplectic FLAVOR

If the phase space is \( \mathbb{R}^d \times \mathbb{R}^d \), then an example of symplectic FLAVOR is obtained from Theorem 2.2 by choosing \( \Phi^\alpha_h \) to be the symplectic Euler (also known as Variational Euler or VE for short) integrator defined by

\[ \Phi^\alpha_h(q, p) = \left( q + h \left( M^{-1} \left( p - h (V(q) + \alpha U(q)) \right) \right), p \right) \]  

(2.9)

and letting \( \Theta^\delta \) be defined by (2.5).
2.1.3 An example of a symplectic and time-reversible FLAVOR

If the phase space is the Euclidean space $\mathbb{R}^d \times \mathbb{R}^d$, then an example of symplectic and time-reversible FLAVOR is obtained by letting $\Theta_\delta$ be defined by Equation (2.7) of Theorem 2.3 by choosing $\Phi_h^{\alpha}$ to be the symplectic Euler integrator defined by (2.9) and

$$
\Phi_h^{\alpha,*}(q,p) = \begin{pmatrix} q \\ p \end{pmatrix} + h \begin{pmatrix} M^{-1}p \\ -V(q + hM^{-1}p) - \alpha U(q + hM^{-1}p) \end{pmatrix}.
$$

(2.10)

2.1.4 An artificial FLAVOR

There is not a unique way of averaging the flows of (2.2). We present below an alternative method based on the freezing and unfreezing of degrees of freedom associated with fast potentials. We have called this method “artificial” because it is intrusive. With this method, the discrete flow approximating solutions of (2.1) is given by (2.4) with

$$
\Theta_\delta := \theta_\delta^r \circ \theta_\tau^e \circ \theta_\delta^V
$$

(2.11)

where $\theta_\delta^V$ is a symplectic map corresponding to the flow of $H_{\text{slow}}(q,p) := V(q)$, approximating the effects of the soft potential on momentum over the mesoscopic time step $\delta$ and defined by

$$
\theta_\delta^V(q,p) = (q, p - \delta \nabla V(q)).
$$

(2.12)

$\theta_\tau^e$ is a symplectic map approximating the flow of $H_{\text{fast}}(q,p) := \frac{1}{2} p^T M^{-1} p + \frac{1}{\epsilon} U(q)$ over a microscopic time step $\tau$:

$$
\theta_\tau^e(q,p) = \left( q + \tau M^{-1} p, p - \frac{\tau}{\epsilon} \nabla U(q + t M^{-1} p) \right).
$$

(2.13)

$\theta_\delta^r$ is a map approximating the flow of the Hamiltonian $H_{\text{free}}(q,p) := \frac{1}{2} p^T M^{-1} p$ under holonomic constraints imposing the freezing of stiff variables. Velocities along the direction of constraints have to be stored and set to be 0 before the constrained dynamics, i.e., frozen, and the stored velocities should be restored after the constrained dynamics, i.e., unfrozen; geometrically speaking, one projects to the constrained sub-symplectic manifold, runs the constrained dynamics, and lifts back to the original full space. Oftentimes, the exact solution to the constrained dynamics can be found (examples given in Subsections 5.3, 5.2, 6.2, 6.3 and 6.4).

When the exact solution to the constrained dynamics cannot be easily found, one may want to employ integrators for constrained dynamics such as SHAKE [91] or RATTLE [4] instead. This has to be done with caution, because symplecticity of the translational flow may be lost. The composition of projection onto the constrained manifold (freezing), evolution on the constrained manifold, and lifting from it to the unconstrained space (unfreezing) preserves symplecticity in the unconstrained space only if the evolution on the constrained manifold preserves the inherited symplectic form. A numerical integration preserves the discrete symplectic form on the constrained manifold, but not necessarily the projected continuous symplectic form.
Remark 2.2. This artificial FLAVOR is locally a perturbation of nonintrusive FLAVORS. By splitting theory [81, 54],

\[ \theta_{\delta - \tau}^{tr} \circ \theta_{\tau}^{V} \approx \theta_{\delta - \tau}^{tr} \circ \theta_{\tau}^{V} \circ \theta_{\tau}^{V} \circ \theta_{\tau}^{V} \circ \Phi_{\tau}^{1} \]

(2.14)

whereas \( \Phi_{\delta - \tau}^{0} \circ \Phi_{\tau}^{1} \approx \theta_{\delta - \tau}^{tr} \circ \theta_{\tau}^{V} \circ \Phi_{\tau}^{1} \), where \( \theta_{\tau}^{V} \) is the flow of \( H_{free}(q, p) \) under no constraint. The only difference is that constraints are treated in \( \theta_{\tau}^{V} \) but not in \( \theta_{\tau}^{V} \).

Remark 2.3. This artificial FLAVOR can be formally regarded as \( \Phi_{\delta - \tau}^{\infty} \circ \Phi_{\tau}^{1} \). In contrast natural FLAVOR is \( \Phi_{\delta - \tau}^{0} \circ \Phi_{\tau}^{1} \).

The advantage of this artificial FLAVOR lies in the fact that only \( \tau \ll \sqrt{\epsilon} \ll \delta \) and \( \delta \ll \frac{\tau^{2}}{\epsilon} \ll \delta \) are required for its accuracy (and not \( \frac{\tau^{2}}{\epsilon} \ll \delta \)). We also observe that, in general, artificial FLAVOR overperforms nonintrusive FLAVOR in FPU long time \( (O(\omega^{2})) \) simulations (we refer to Subsection 6.3).

2.2 Variational derivation of FLAVORS

FLAVORS based on variational legacy integrators [78] are variational too. Recall that discrete Lagrangian \( L_{d} \) is an approximation of the integral of the continuous Lagrangian over one time step, and Discrete Euler-Lagrangian equation (DEL) corresponds to the critical point of the discrete action, which is a sum of the approximated integrals. The following diagram commutes:

\[
\begin{array}{ccc}
\text{Singlescale } L_{d} & \xrightarrow{\text{FLAVORization}} & \text{Multiscale } L_{d} \\
\downarrow \text{action principle} & & \downarrow \text{action principle} \\
\text{Singlescale DEL} & \xrightarrow{\text{FLAVORization}} & \text{Multiscale DEL}
\end{array}
\]

For example, recall Variational Euler (i.e. symplectic Euler) for system (2.2) with time step \( h \)

\[
\begin{align*}
p_{k+1} &= p_{k} - h[\nabla V(q_{k}) + \frac{1}{\epsilon} \nabla U(q_{k})] \\
q_{k+1} &= q_{k} + hp_{k+1}
\end{align*}
\]

(2.15)

can be obtained by applying variational principle to the following discrete Lagrangian

\[
L_{dh}^{1/\epsilon}(q_{k}, q_{k+1}) = h \left[ \frac{1}{2} \left( \frac{q_{k+1} - q_{k}}{h} \right)^{2} - \left( V(q_{k}) + \frac{1}{\epsilon} U(q_{k}) \right) \right].
\]

(2.16)

Meanwhile, FLAVORized Variational Euler with smallstep \( \tau \) and mesostep \( \delta \)

\[
\begin{align*}
p'_{k} &= p_{k} - \tau[\nabla V(q_{k}) + \frac{1}{\tau} \nabla U(q_{k})] \\
q'_{k} &= q_{k} + \tau p'_{k} \\
p_{k+1} &= p'_{k} - (\delta - \tau) \nabla V(q'_{k}) \\
q_{k+1} &= q'_{k} + (\delta - \tau)p_{k+1}
\end{align*}
\]

(2.17)
can be obtained by applying variational principle to the FLAVORized discrete Lagrangian
\[ L_{d\delta}(q_k, q'_k, q_{k+1}) = L_{d\tau}^{1/\epsilon}(q_k, q'_k) + L_{d\delta-\tau}(q'_k, q_{k+1}) \]
\[ = \tau \left[ \frac{1}{2} \left( \frac{q'_k - q_k}{\tau} \right)^2 - \left( V(q_k) + \frac{1}{\epsilon} U(q_k) \right) \right] + (\delta - \tau) \left[ \frac{1}{2} \left( \frac{q_{k+1} - q'_k}{\delta - \tau} \right)^2 - V(q'_k) \right] \]
(2.18)

FLAVORizations of other variational integrators such as Velocity Verlet follow similarly.

3 SDEs

Asymptotic problems for stochastic differential equations arose and were solved simultaneously with the very beginnings of the theory of such equations [104]. Here, we refer to the early work of Gikhman [49], Krylov [67, 68], Bogolyubov [13] and Papanicolaou-Kohler [87]. We refer in particular to Skorokhod’s detailed monograph [104]. As for ODEs, effective equations for stiff SDEs can be obtained by averaging the instantaneous coefficients (drift and the diffusivity matrix squared) with respect to the fast components; we refer to Chapter II, Section 3 of [104] for a detailed analysis including error bounds. Numerical methods such as HMM [37] and equation-free methods [7] have been extended to SDEs based on this averaging principle. Implicit methods in general fail to capture the effective dynamics of the slow time scale because they cannot correctly capture non-Dirac invariant distributions [76] (we refer to non-Dirac invariant distribution as a measure of probability on the configuration space whose support is not limited to a single point). Another idea is to treat fast variables by conditioning; here, we refer to optimal prediction [28, 27, 29] that has also been used for model reduction. We also refer to [8, 52, 108, 22, 23, 76, 2].

Since FLAVORS are obtained via flow averaging, they have a natural extension to SDEs developed in this section. As for ODEs, FLAVORS are directly applied to SDEs with mixed (hidden) slow and fast variables without prior (analytical or numerical) identification of slow variables. Furthermore, they can be implemented using a pre-existing scheme by turning on and off the stiff parameters.

For the sake of clarity, we will start the description of with the following SDE on \( \mathbb{R}^d \):
\[ du_t' = \left( G(u_t') + \frac{1}{\epsilon} F(u_t') \right) \, dt + \left( H(u_t') + \frac{1}{\sqrt{\epsilon}} K(u_t') \right) \, dW_t, \quad u_0' = u_0 \]
(3.1)
where \((W_t)_{t \geq 0}\) is a \( d \)-dimensional Brownian Motion; \( F \) and \( G \) are vector fields on \( \mathbb{R}^d \); \( H \) and \( K \) are \( d \times d \) matrix fields on \( \mathbb{R}^d \). In Subsection [3.5] we will consider the more general form [3.15].

**Condition 3.1.** Assume that:

1. \( F, G, H \) and \( K \) are uniformly bounded and Lipschitz continuous.
Remark 3.1 Two-scale flow convergence for SDEs

towards infinity. To simplify the presentation, we will use the global uniform regularity. can be relaxed to local regularity by adding a control on the rate of escape of the process measurable) stochastic process on \( \mathbb{R} \) continuous on \( \mathbb{R} \), and for all \( t > 0 \)

\[
\begin{align*}
\frac{1}{T} \int_0^T \mathbb{E} \left[ \phi(Y_s) \right] - \int_0^T \phi(y) \mathbb{E} \left[ \mu(x_0, dy) \right] \leq \chi \left( \left\| (x_0, y_0) \right\| \right) \mathbb{E} \left[ \max_{0 \leq t \leq 3} \left\| \phi \right\|_{C^r} \right]
\end{align*}
\]

where \( r \mapsto \chi(r) \) is bounded on compact sets.

4. For all \( u_0, T > 0 \), \( \sup_{0 \leq t \leq T} \mathbb{E} \left[ \chi \left( \left\| u_t \right\| \right) \right] \) is uniformly bounded in \( \epsilon \).

Remark 3.1. As in the proof of Theorem 1.1 the uniform regularity of \( F, G, H \) and \( K \) can be relaxed to local regularity by adding a control on the rate of escape of the process towards infinity. To simplify the presentation, we will use the global uniform regularity.

We will now extend the definition of two-scale flow convergence introduced in Subsection 1.2 to stochastic processes.

3.1 Two-scale flow convergence for SDEs

Let \( (\xi^\epsilon(t, \omega))_{t \in \mathbb{R}^+} \) be a sequence of stochastic processes on \( \mathbb{R}^d \) (progressively measurable mappings from \( \mathbb{R}^+ \times \Omega \) to \( \mathbb{R}^d \)) indexed by \( \epsilon > 0 \). Let \( (X_t)_{t \in \mathbb{R}^+} \) be a (progressively measurable) stochastic process on \( \mathbb{R}^{d-p} \) \( (p \geq 0) \). Let \( x \mapsto \nu(x, dz) \) be a function from \( \mathbb{R}^{d-p} \) into the space of probability measures on \( \mathbb{R}^d \).

Definition 3.1. We say that the process \( \xi^\epsilon \) \( F \)-converges to \( \nu(X_t, dz) \) as \( \epsilon \downarrow 0 \) and write \( \xi^\epsilon \xrightarrow{F_{\epsilon \downarrow 0}} \nu(X_t, dz) \) if and only if for all function \( \varphi \) bounded and uniformly Lipshitzcontinuous on \( \mathbb{R}^d \), and for all \( t > 0 \)

\[
\lim_{h \to 0} \lim_{\epsilon \to 0} \frac{1}{h} \int_t^{t+h} \mathbb{E} \left[ \varphi(\xi^\epsilon_s) \right] ds = \mathbb{E} \left[ \int_{\mathbb{R}^d} \varphi(z) \nu(X_t, dz) \right]
\]
3.2 Non intrusive FLAVORS for SDEs

Let $\omega$ be a random sample from a probability space $(\Omega, \mathcal{F}, P)$ and $\Phi^\alpha_h(\cdot, \omega)$ a random mapping from $\mathbb{R}^d$ onto $\mathbb{R}^d$ approximating the flow of (3.1) for $\alpha = 1/\epsilon$. If the parameter $\alpha$ can be controlled, then $\Phi^\alpha_h$ can be used as a black box for accelerating the computation of solutions of (3.1) without prior identification of slow variables. Indeed, assume that there exists a constant $h_0 > 0$ and a normal random vector $\xi(\omega)$ such that for $h \leq h_0 \min(\frac{1}{\alpha}, 1)$

$$\left(\mathbb{E}\left[ |\Phi^\alpha_h(u, \omega) - u - hG(u) - \sqrt{hH(u)\xi(\omega)} - \sqrt{\alpha h F(u)} - \sqrt{\alpha h K(u)}\xi(\omega)|^2 \right] \right)^{\frac{1}{2}} \leq Ch^\frac{3}{2}(1+\alpha)^{\frac{3}{2}} \tag{3.6}$$

then FLAVOR can be defined as the algorithm simulating the stochastic process

$$\begin{cases}
\bar{u}_0 = u_0 \\
\bar{u}_{(k+1)\delta} = \Phi^0_{\delta - \tau}(\cdot, \omega'_k) \circ \Phi^\frac{1}{\epsilon}(\bar{u}_{k\delta}, \omega_k) \\
\bar{u}_t = \bar{u}_{k\delta} \text{ for } k\delta \leq t < (k+1)\delta
\end{cases} \tag{3.7}$$

where $\omega_k, \omega'_k$ are i.i.d. samples from the probability space $(\Omega, \mathcal{F}, P)$, $\delta \leq h_0$ and $\tau \in (0, \delta)$ such that $\tau \leq \tau_0 \epsilon$. Theorem 3.1 establishes the asymptotic accuracy of FLAVOR for $\tau \ll \epsilon \ll \delta$ and

$$\left(\frac{\tau}{\epsilon}\right)^2 \ll \delta \ll \frac{\tau}{\epsilon}. \tag{3.8}$$

3.3 Convergence theorem

**Theorem 3.1.** Let $u^\epsilon$ be the solution to (3.1) and $\bar{u}_t$ defined by (3.7). Assume that equation (3.6) and Conditions 3.1 are satisfied, then

- $u^\epsilon_t$ $F$-converges towards $\eta^{-1} \ast (\delta X_t \otimes \mu(X_t, dy))$ as $\epsilon \downarrow 0$ where $X_t$ is the solution to

$$dX_t = \int g(X_t, y) \mu(X_t, dy) \, dt + \bar{\sigma}(X_t) \, dB_t \quad X_0 = x_0 \tag{3.9}$$

where $\bar{\sigma}$ is a $(d - p) \times (d - p)$ matrix field defined by

$$\bar{\sigma}\bar{\sigma}^T = \int \sigma\sigma^T(x, y) \mu(x, dy) \tag{3.10}$$

and $B_t$ a $(d - p)$-dimensional Brownian Motion.

- $\bar{u}_t$ $F$-converges towards $\eta^{-1} \ast (\delta X_t \otimes \mu(X_t, dy))$ as $\epsilon \downarrow 0$, $\tau \leq \delta$, $\frac{\tau}{\epsilon} \downarrow 0$, $\frac{\delta}{\tau} \downarrow 0$ and

$$\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}} \frac{1}{\delta} \downarrow 0.$$

The proof of convergence of SDEs of type (3.2) is classical, and a comprehensive monograph can be found in Chapter II of [104]. A proof of (mean squared) convergence of HMM applied to (3.2) (separated slow and fast variables) with $\sigma = 0$ has been
obtained in [37]. A proof of (mean squared) convergence of the Equation-Free Method applied to (3.2) with \( \sigma \neq 0 \) but independent of fast variables has been obtained in [50]. Theorem 3.1 proves the convergence in distribution of FLAVOR applied to SDE (3.1) with hidden slow and fast processes. One of the main difficulties of the proof of Theorem 3.1 lies in the fact that we are not assuming that the noise on (hidden) slow variables is null or independent from fast variables. Without this assumption, \( x^\epsilon_t \) converges only weakly towards \( X_t \), the convergence of \( u^\epsilon \) can only be weak and techniques for strong convergence can not be used. The proof of Theorem 3.1 relies on a powerful result by Skorokhod (Theorem 1 of Chapter II of [104]) stating that the convergence in distribution of a sequence of stochastic processes is implied by the convergence of their generators. We refer to Subsection 7.2 of the appendix for the detailed proof of Theorem 3.1.

3.4 Natural FLAVORS

As for ODEs, it is not necessary to use legacy integrators to obtain FLAVORS for SDEs. More precisely, Theorem 3.1 remains valid if FLAVORS are defined to be algorithms simulating the discrete process

\[
\begin{align*}
\bar{u}_0 &= u_0 \\
\bar{u}_{(k+1)\delta} &= \theta^G_{\delta-\tau}(\cdot, \omega^\epsilon_k) \circ \theta^\epsilon_{\delta}(\bar{u}_{k\delta}, \omega_k) \\
\bar{u}_t &= \bar{u}_{k\delta} \quad \text{for} \quad k\delta \leq t < (k+1)\delta
\end{align*}
\tag{3.11}
\]

where \( \omega_k, \omega^\epsilon_k \) are i.i.d. samples from the probability space \((\Omega, \mathcal{F}, \mathbb{P})\) and \( \theta^\epsilon_{\delta} \) and \( \theta^G_{\delta-\tau} \) are two random mappings from \( \mathbb{R}^d \) onto \( \mathbb{R}^d \) satisfying following conditions 3.2. More precisely, \( \theta^\epsilon_{\delta}(\cdot, \omega) \) approximates in distribution the flow of (3.1) over time steps \( \tau \ll \epsilon \). \( \theta^G_{\delta}(\cdot, \omega) \) approximates in distribution the flow of

\[
dv_t^\epsilon = G(v^\epsilon_t) dt + H(v^\epsilon_t) dW_t
\tag{3.12}
\]

over time steps \( h \ll 1 \).

Condition 3.2. Assume that:

1. There exists \( h_0, C > 0 \) and a \( d \)-dimensional centered Gaussian vector \( \xi(\omega) \) with identity covariance matrix such that for \( h \leq h_0 \),

\[
\left( \mathbb{E} \left[ \left| \theta^G_h(u, \omega) - u - hG(u) - \sqrt{h}H(u)\xi(\omega) \right|^2 \right] \right)^{\frac{1}{2}} \leq Ch^{\frac{3}{2}}
\tag{3.13}
\]

2. There exists \( \tau_0, C > 0 \) and a \( d \)-dimensional centered Gaussian vector \( \xi(\omega) \) with identity covariance matrix such that for \( \frac{\tau}{\epsilon} \leq \tau_0 \),

\[
\left( \mathbb{E} \left[ \left| \theta^\epsilon_{\tau}(u, \omega) - u - \tau G(u) - \frac{\tau}{\epsilon}F(u) - \sqrt{\tau}H(u)\xi(\omega) - \sqrt{\tau}K(u)\xi(\omega) \right|^2 \right] \right)^{\frac{1}{2}} \leq C\left( \frac{\tau}{\epsilon} \right)^{\frac{3}{2}}
\tag{3.14}
\]
3. For all \( u_0, T > 0 \), \( \sup_{0 \leq n \leq T/\delta} E \left[ \chi \left( \|u_{n\delta}\| \right) \right] \) is uniformly bounded in \( \epsilon, 0 < \delta \leq h_0, \tau \leq \min(\tau_0, \delta) \), where \( \hat{u} \) is defined by (3.11).

3.5 FLAVORS for generic stiff SDEs

FLAVORS for stochastic systems have a natural generalization to SDEs on \( \mathbb{R}^d \) of the form

\[
du^{\alpha,\epsilon} = F(u^{\alpha,\epsilon}, \alpha, \epsilon) \, dt + K(u^{\alpha,\epsilon}, \alpha, \epsilon) \, dW_t
\]

(3.15)

where \( (W_t)_{t \geq 0} \) is a \( d \)-dimensional Brownian Motion, \( F \) and \( K \) are Lipshitz continuous in \( u \).

**Condition 3.3.** Assume that:

1. \( \gamma \mapsto F(u, \alpha, \gamma) \) and \( \gamma \mapsto K(u, \alpha, \gamma) \) are uniformly continuous in the neighborhood of 0.

2. There exists a diffeomorphism \( \eta := (\eta^x, \eta^y) \), from \( \mathbb{R}^d \) onto \( \mathbb{R}^{d-p} \times \mathbb{R}^p \), independent from \( \epsilon, \alpha \), with uniformly bounded \( C^1, C^2 \) and \( C^3 \) derivatives, and such that the stochastic process \((x_t^\alpha, y_t^\alpha) = (\eta^x(u_t^{\alpha,0}), \eta^y(u_t^{\alpha,0}))\) satisfies for all \( \alpha \geq 1 \) the SDE

\[
dx^\alpha = g(x^\alpha, y^\alpha) \, dt + \sigma(x^\alpha, y^\alpha) \, dW_t \quad x_0^\alpha = x_0
\]

(3.16)

where \( g \) is \( d - p \) dimensional vector field, \( \sigma \) is a \((d - p) \times d\)-dimensional matrix field, \( g \) and \( \sigma \) are uniformly bounded and Lipschitz continuous in \( x \) and \( y \).

3. There exists a family of probability measures \( \mu(x, dy) \) on \( \mathbb{R}^p \) such that for all \( x_0, y_0, T \) \((x_0, y_0) := \eta(u_0)\) and \( \varphi \) with uniformly bounded \( C^r \) derivatives for \( r \leq 3 \),

\[
\frac{1}{T} \int_0^T \mathbb{E} [\varphi(y^n)] \, ds - \int \varphi(y) \mu(x_0, dy) \leq \chi(\|x_0, y_0\|) \left( E_1(T) + E_2(T^{3/2}) \right) \max_{r \leq 3} \|\varphi\|_{C^r}
\]

(3.17)

where \( r \mapsto \chi(r) \) is bounded on compact sets and \( E_2(r) \to 0 \) as \( r \to \infty \) and \( E_1(r) \to 0 \) as \( r \to 0 \).

4. For all \( u_0, T > 0 \), \( \sup_{0 \leq t \leq T} \mathbb{E} \left[ \chi \left( \|u_{t}^{0,0}\| \right) \right] \) is uniformly bounded in \( \alpha \geq 1 \).

**Remark 3.2.** As in the proof of Theorem 1.1, the uniform regularity of \( g \) and \( \sigma \) can be relaxed to local regularity by adding a control on the rate of escape of the process towards infinity. To simplify the presentation, we have use the global uniform regularity.

Let \( \omega \) be a random sample from a probability space \((\Omega, \mathcal{F}, \mathbb{P})\) and \( \Phi_{h,\epsilon}^{\alpha,\epsilon}(., \omega) \) a random mapping from \( \mathbb{R}^d \) onto \( \mathbb{R}^d \) approximating in distribution the flow of (3.15) over time steps \( \tau \ll \epsilon \). If the parameter \( \alpha \) can be controlled, then \( \Phi_{h,\epsilon}^{\alpha,\epsilon} \) can be used as a black box for accelerating the computation of solutions of (3.15). The acceleration is obtained without prior identification of the slow variables.

**Condition 3.4.** Assume that:
1. There exists $h_0, C > 0$ and a $d$-dimensional centered Gaussian vector $\xi(\omega)$ with identity covariance matrix such that for $h \leq h_0$, $0 < \epsilon \leq 1 \leq \alpha$ and $h \leq h_0 \min(\frac{1}{\alpha^\nu}, 1)$

$$\left( \mathbb{E} \left[ \Phi_h^\alpha(\epsilon)(u) - u - hF(u, \alpha, \epsilon) - \sqrt{h}\xi(\omega)K(u, \alpha, \epsilon) \right]^2 \right)^{\frac{1}{2}} \leq Ch^\frac{3}{2}(1 + \alpha^\nu) \quad (3.18)$$

2. For all $u_0$, $T > 0$, $\sup_{0 \leq n \leq T/h} \mathbb{E} \left[ \chi\left(\|\bar{\eta}_n\|\right) \right]$ is uniformly bounded in $\epsilon$, $0 < \delta \leq h_0$, $\tau \leq \min(h_0\epsilon^\nu, \delta)$, where $\bar{u}$ is defined by (3.19).

**FLAVORS** Let $\delta \leq h_0$ and $\tau \in (0, \delta)$ such that $\tau \leq \tau_0\epsilon^\nu$. We define FLAVORS as the class of algorithms simulating the stochastic process $t \mapsto \bar{u}_t$ defined by

$$\begin{cases}
\bar{u}_0 = u_0 \\
\bar{u}_{(k+1)\delta} = \Phi^\alpha_{\delta-\tau}(., \omega'_k) \circ \Phi^\frac{1}{2}_{\delta-\tau}(\bar{u}_k, \omega_k) \\
\bar{u}_t = \bar{u}_{k\delta} \text{ for } k\delta \leq t < (k+1)\delta
\end{cases} \quad (3.19)$$

where $\omega_k, \omega'_k$ are i.i.d. samples from the probability space $(\Omega, \mathcal{F}, P)$.

**Remark 3.3.** $\omega_k$ simulates the randomness of the increment of the Brownian Motion between times $\delta k$ and $\delta k + \tau$. $\omega'_k$ simulates the randomness of the increment of the Brownian Motion between times $\delta k + \tau$ and $\delta(k+1)$. The independence of $\omega_k$ and $\omega'_k$ is reflection of the independence of the increments of a Brownian Motion.

The following theorem shows that the flow averaging integrator is accurate with respect to $F$-convergence for $\tau \ll \epsilon^\nu \ll \delta$ and

$$\left(\frac{\tau}{\epsilon^\nu}\right)^\frac{3}{2} \ll \delta \ll \frac{\tau}{\epsilon^\nu}. \quad (3.20)$$

**Theorem 3.2.** Let $u_t^\frac{1}{\epsilon^\nu}$ be the solution to (3.15) with $\alpha = \frac{1}{\epsilon}$ and $\bar{u}_t$ be defined by (3.19). Assume that Conditions 3.3 and 3.4 are satisfied then

- $u_t^\frac{1}{\epsilon^\nu} F$-converges towards $\eta^{-1} \ast (\delta X_t \otimes \mu(X_t, dy))$ as $\epsilon \downarrow 0$ where $X_t$ is the solution to

$$dX_t = \int g(X_t, y) \mu(X_t, dy) + \tilde{\sigma}(X_t) dB_t \quad X_0 = x_0 \quad (3.21)$$

where $\tilde{\sigma}$ is a $(d-p) \times (d-p)$ matrix field defined by

$$\tilde{\sigma}\tilde{\sigma}^T = \int \sigma\sigma^T(x, y) \mu(x, dy) \quad (3.22)$$

and $B_t$ a $(d-p)$-dimensional Brownian Motion.

- As $\epsilon \downarrow 0$, $\tau \epsilon^{-\nu} \downarrow 0$, $\delta \epsilon^{-\nu} \downarrow 0$, $\left(\frac{\tau}{\delta}\right)^\frac{3}{2} \downarrow 0$, $\bar{u}_t F$-converges towards $\eta^{-1} \ast (\delta X_t \otimes \mu(X_t, dy))$ as $\epsilon \downarrow 0$ where $X_t$ is the solution to (3.21).
Proof. The proof of Theorem 3.2 is similar to the proof of Theorem 3.1. The condition \( \epsilon \ll 1 \) is needed for the approximation of \( u^{\alpha,\epsilon} \) by \( u^{\alpha,0} \) and for the \( F \)-convergence of \( u^{1/\epsilon} \).

Since \( y_t^\alpha = \eta^\alpha(u_t^\alpha,0) \) the condition \( \tau \ll \epsilon \nu \) is used along with Equation (3.18) for the accuracy of \( \Phi_{1/\epsilon}^{1/\epsilon} \) in (locally) approximating \( y^\alpha_t \). The condition \( \delta \ll \tau \epsilon \nu \) allows for the averaging of \( g \) and \( \sigma \) to take place prior to a significant change of \( x^\alpha_t \); more precisely, it allows for \( m \gg 1 \) iterations of \( \Phi_{1/\epsilon}^{1/\epsilon} \) prior to a significant change of \( x^\alpha_t \). The condition \( (\tau \epsilon \nu)^{3/2} \ll \delta \) is required in order to control the error accumulated by \( m \) iterations of \( \Phi_{1/\epsilon}^{1/\epsilon} \).

□

4 Stochastic mechanical systems: Langevin equations

Since the foundational work of Bismut [12], the field of stochastic geometric mechanics has grown in response to the demand for tools to analyze the structure of continuous and discrete mechanical systems with uncertainty [103, 57, 112, 30, 82, 83, 85, 69, 77, 18, 17, 19]. Like their deterministic counterparts, these integrators are structure preserving in terms of statistical invariants.

In this section, FLAVORS are developed to be structure preserving integrators for stiff stochastic mechanical systems, i.e., stiff Langevin equations of the form

\[
\begin{align*}
\dot{q} &= M^{-1}p \\
\dot{p} &= -\nabla V(q) dt - \frac{1}{\epsilon} \nabla U(q) dt - cp dt + \sqrt{2\beta^{-1}c^2} dW_t
\end{align*}
\] (4.1)

and of the form

\[
\begin{align*}
\dot{q} &= M^{-1}p \\
\dot{p} &= -\nabla V(q) dt - \frac{1}{\epsilon} \nabla U(q) dt - c\epsilon p dt + \sqrt{2\beta^{-1}c^2} \frac{1}{\sqrt{\epsilon}} dW_t
\end{align*}
\] (4.2)

where \( c \) is a positive symmetric \( d \times d \) matrix.

Remark 4.1. Provided that hidden fast variables remain locally ergodic, one can also consider Hamiltonians with a mixture of both slow and fast noise and friction. For the sake of clarity, we have restricted our presentation to (4.1) and (4.2).

Equations (4.1) and (4.2) model a mechanical system with Hamiltonian

\[
H(q,p) := \frac{1}{2}p^T M^{-1} p + V(q) + \frac{1}{\epsilon} U(q).
\] (4.3)

The phase space is the Euclidean space \( \mathbb{R}^d \times \mathbb{R}^d \) or a cotangent bundle \( T^*\mathcal{M} \) of a configuration manifold \( \mathcal{M} \).

Remark 4.2. If \( c \) is not constant and \( \mathcal{M} \) is not the usual \( \mathbb{R}^d \times \mathbb{R}^d \) Euclidean space, one should use the Stratonovich integral instead of the Itô integral.
4.1 FLAVORS for stochastic mechanical systems on manifolds

As in Section 2, we assume that we are given a mapping \( \Phi^\alpha_h \) acting on the phase space such that for \( h \leq h_0 \min(1, \alpha^{-\frac{1}{2}}) \)

\[
\left| \Phi^\alpha_h(q,p) - (q,p) - h \left( M^{-1} p, -V(q) - \alpha U(q) \right) \right| \leq Ch^2(1 + \alpha)
\]  

(4.4)

Next, consider the following Ornstein-Uhlenbeck equations:

\[
dp = -\alpha c p \, dt + \sqrt{\alpha} \sqrt{2} \beta^{-1} c^{1/2} dW_t
\]  

(4.5)

The stochastic flow of (4.5) is defined by the following stochastic evolution map:

\[
\Psi_{t_{1,1},t_{1,2}}^\alpha(q,p) = \left( q, e^{-\alpha \left( t_2 - t_1 \right)} p + \sqrt{2} \beta^{-1} c \int_{t_1}^{t_2} e^{-\alpha \left( t - t_1 \right)} dW_s \right)
\]  

(4.6)

Let \( \delta \leq h_0 \) and \( \tau \in (0, \delta) \) such that \( \tau \leq \tau_0 / \sqrt{\alpha} \). FLAVOR for (4.1) can then be defined by

\[
\begin{align*}
\bar{q}_0 & = (q_0, p_0) \\
\bar{q}_{(k+1)\delta} & = \Phi^\delta_{t_{1,1},t_{1,2}} \circ \Psi_{t_{1,1},t_{1,2}}^\delta \circ \Phi^\frac{1}{\tau}_{t_{1,2},t_{1,2}} \circ \Psi_{t_{1,2},t_{1,2}}^\frac{1}{\tau} \circ \Phi^\frac{1}{\tau}_{t_{1,2},t_{1,2}}(q_{k\delta}, p_{k\delta})
\end{align*}
\]  

(4.7)

and FLAVOR for (4.2) can be defined by

\[
\begin{align*}
\bar{q}_0 & = (q_0, p_0) \\
\bar{q}_{(k+1)\delta} & = \Phi^\delta_{t_{1,1},t_{1,2}} \circ \Phi^\frac{1}{\tau}_{t_{1,2},t_{1,2}} \circ \Psi_{t_{1,2},t_{1,2}}^\frac{1}{\tau} \circ \Phi^\frac{1}{\tau}_{t_{1,2},t_{1,2}}(q_{k\delta}, p_{k\delta})
\end{align*}
\]  

(4.8)

Theorem 3.2 establishes the accuracy of these integrators under Conditions 3.3 and 3.4 provided that \( \tau \ll \sqrt{\epsilon} \ll \delta \) and \( \left( \frac{\tau}{\sqrt{\epsilon}} \right)^{3/2} \ll \delta \ll \frac{\tau}{\sqrt{\epsilon}} \).

4.2 Structure Preserving properties of FLAVORS for stochastic mechanical systems on manifolds

First, observe that if \( \Phi^\alpha_h \) and \( \Psi^\frac{1}{\tau}_h \) are symmetric under a group action for all \( \epsilon > 0 \), then the resulting FLAVOR, as a symmetric composition of symmetric steps, is symmetric under the same group action (see comment below Theorem 2.3).

Similarly, the following theorem shows that FLAVORS inherits structure-preserving properties from those associated with \( \Phi^\alpha_h \) (the component approximating the Hamiltonian part of the flow).

Theorem 4.1.

- If \( \Phi^\alpha_h \) is symplectic, then the FLAVORS defined by (4.7) and (4.8) are quasi-symplectic as defined in Conditions RL1 and RL2 of [84] (it degenerates to a symplectic method if friction is set equal to zero and the Jacobian of the flow map is independent of \( (q,p) \)).
• If in addition $c$ is isotropic then FLAVOR defined by (4.7) is conformally symplectic, i.e., it preserves the precise symplectic area change associated to the flow of inertial Langevin processes \[80\].

Proof. Those properties are a consequence of the fact that FLAVORS are splitting schemes. The quasi-symplecticity and symplectic conformallity of GLA has been obtained in a similar way in \[17\].

4.2.1 Example of quasi-symplectic FLAVORS

An example of quasi-symplectic FLAVOR can be obtained by choosing $\Phi^c_h$ to be the symplectic Euler integrator defined by (2.9). This integrator is also conformally symplectic if $c$ is isotropic and friction is slow.

4.2.2 Example of quasi-symplectic and time-reversible FLAVORS

Defining $\Phi^c_h$ by (2.9) and $\Phi^c_{h,*}$ by (2.10), an example of quasi-symplectic and time-reversible FLAVOR can be obtained by using the symmetric Strang splitting:

\[
(\bar{q}(k+1),\bar{p}(k+1)) = \Psi^{1}_{k\delta+k\delta} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Psi^{1}_{k\delta+k\delta} (q,p) \tag{4.9}
\]

for (4.1) and

\[
(\bar{q}(k+1),\bar{p}(k+1)) = \Psi^{1}_{(k+1)\delta} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Phi_{\frac{1}{2}} \circ \Psi^{1}_{(k+1)\delta+k\delta} (q,p) \tag{4.10}
\]

for (4.2). This integrator is also conformally symplectic if $c$ is isotropic and friction is slow.

4.2.3 Example of Boltzmann-Gibbs reversible Metropolis-adjusted FLAVORS

Since the probability density of $\Psi_{t_{1},t_{2}}$ can be explicitly computed, it follows that the probability densities of (4.9) and (4.10) can be explicitly computed, and these algorithms can be metropolized and made reversible with respect to the Gibbs distribution as it has been shown in \[19\] for the Geometric Langevin Algorithm introduced in \[17\]. This metropolization leads to stochastically stable (and ergodic if the noise applied on momentum is not degenerate) algorithms. We refer to \[19\] for details. Observe that if the proposed move is rejected, the momentum has to be flipped and the acceptance probability involves a momentum flip. It is proven in \[19\] that GLA \[17\] remains strongly accurate after a metropolization involving local momentum flips. Whether this preservation of accuracy over trajectories transfers in a weak sense (in distributions) to FLAVORS remains to be investigated.
5 Numerical analysis of FLAVOR based on Variational Euler

5.1 Stability

Consider the following linear Hamiltonian system

\[ H(x, y, p_x, p_y) = \frac{1}{2}p_x^2 + \frac{1}{2}p_y^2 + \frac{1}{2}x^2 + \frac{\omega^2}{2}(y - x)^2 \]  

(5.1)

with \( \omega \gg 1 \). Here \( \frac{x+y}{2} \) is the slow variable and \( y - x \) is the fast variable.

It can be shown that, when applied to (5.1), Symplectic Euler (2.9) is stable if and only if \( h \leq \sqrt{2}/\omega \). Write \( \Theta_{\delta,\tau}^{\delta,\tau} \) the non-intrusive FLAVOR (2.5) obtained by using Symplectic Euler (2.9) as a Legacy integrator. Write \( \Theta_{\delta,\tau}^{\delta,\tau} \) the artificial FLAVOR described in Subsection 2.1.4.

**Theorem 5.1.** The non-intrusive FLAVOR \( \Theta_{\delta,\tau}^{\delta,\tau} \) with \( 1/\sqrt{\tau} \gg \omega \gg 1 \) is stable if and only if \( \delta \in (0, 2) \).

The artificial FLAVOR \( \Theta_{\delta,\tau}^{\delta,\tau} \) with \( 1/\tau \gg \omega \gg 1 \) is stable if and only if \( \delta \in (0, 2\sqrt{2}) \).

**Proof.** The numerical scheme associated with \( \Theta_{\delta,\tau}^{\delta,\tau} \) can be written as

\[
\begin{bmatrix}
  y_{n+1} \\
  x_{n+1} \\
  (p_y)_{n+1} \\
  (p_x)_{n+1}
\end{bmatrix} = T
\begin{bmatrix}
  y_n \\
  x_n \\
  (p_y)_n \\
  (p_x)_n
\end{bmatrix}
\]  

(5.2)
The characteristic polynomial of $T$ is
\[
\lambda^4 + (-4 + \delta^2 - \delta^2 \tau^2 + 2\delta \tau^3 - \tau^4 + 2\delta \tau \omega^2 - \delta^2 \tau^2 \omega^2 + 2\delta \tau^2 \omega^2 - \tau^4 \omega^2)\lambda^3 + (6 - 2\delta^2 + 2\delta^2 \tau^2 - 4\delta \tau^3 + 2\tau^4 - 4\delta \tau \omega^2 + 3\delta^3 \tau \omega^2 + 2\delta^2 \tau^2 \omega^2 - 4\delta^3 \tau^2 \omega^2 - \delta^3 \tau^3 \omega^2 + 2\tau^4 \omega^2 + 2\delta^2 \tau^3 \omega^2 - \tau^4 \delta \tau \omega^2 - 2\delta \tau^3 \omega^2 + 2\delta \tau^3 \omega^2 - \tau^4 \delta^2 \omega^2)\lambda + 1
\] (5.3)

Since $\omega \gg 1$, $\tau \ll 1/\omega^2$, as long as $\delta \lesssim 1$ roots to the above polynomial are close to roots to the asymptotic polynomial
\[
\lambda^4 + (\delta^2 - 4)\lambda^3 + (6 - 2\delta^2)\lambda^2 + (\delta^2 - 4)\lambda + 1
\] (5.4)

which can be shown to be 1 with multiplicity 2 and $\frac{1}{2}(2 - \delta^2 \pm \delta \sqrt{\delta^2 - 4})$. It is easy to see that all roots are complex numbers with moduli less or equal to one if and only if $|\delta| \leq 2$.

The numerical scheme associated with $\Theta^a_{\delta, \tau}$ can be written as in (5.2) with
\[
T = \begin{bmatrix}
1 & 0 & \delta - \tau & 0 \\
0 & 1 & 0 & \delta - \tau \\
0 & 0 & 1 & 0 \\
\end{bmatrix}
\begin{bmatrix}
1 & 0 & 0 & 0 \\
1 & 0 & \tau & 0 \\
-\tau & 0 & 1 & 0 \\
\tau & 0 & 0 & 1 \\
\end{bmatrix}
\begin{bmatrix}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & \tau \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
\end{bmatrix}
\] (5.5)

The characteristic polynomial of $T$ is
\[
2\lambda^3 + (4\omega^2 \tau^2 + \tau \delta + \delta^2 - 8)\lambda^2 + (12 - 2\delta^2 - 2\delta \tau - 8\tau^2 \omega^2 + 2\delta^2 \tau^2 \omega^2)\lambda + (4\omega^2 \tau^2 + \tau \delta + \delta^2 - 8)\lambda + 2
\] (5.6)

Since $\omega \gg 1$, $\tau \ll 1/\omega$, as long as $\delta \lesssim 1$ roots to the above polynomial are close to roots to the asymptotic polynomial
\[
2\lambda^4 + (\delta^2 - 8)\lambda^3 + (12 - 2\delta^2)\lambda^2 + (\delta^2 - 8)\lambda + 1
\] (5.7)

which can be shown to be 1 with multiplicity 2 and $\frac{1}{4}(4 - \delta^2 \pm \delta \sqrt{\delta^2 - 8})$. All roots are complex numbers with moduli less or equal to one if and only if $|\delta| \leq 2\sqrt{2}$.

Figures 2(a) and 2(b) illustrate the domain of stability of nonintrusive FLAVOR (based on symplectic Euler (2.5) and (2.9)) and artificial FLAVOR (2.11) applied to the flow of (5.1), i.e. values of $\delta$ and $\tau/\epsilon$ ensuring stable numerical integrations. We observe that artificial FLAVOR has a much larger stability domain than nonintrusive FLAVOR. Specifically, for nonintrusive FLAVOR and large values of $\delta$, $\tau = o(\sqrt{\tau})$ is not enough.
and one needs $\tau = o(\epsilon)$ for a stable integration, whereas artificial FLAVOR only requires $\tau = \sqrt{2}\epsilon$, a minimum requirement for a stable symplectic Euler integration of the fast dynamics.

Notice that there is no resonance behavior in terms of stability; everything below the two curves is stable and everything outside is not stable (plots not shown).

5.2 Error analysis

The flow of (5.1) has been explicitly computed and compared with solutions obtained from nonintrusive FLAVOR based on symplectic Euler ((2.5) and (2.9)) and with artificial FLAVOR (2.11).

The total simulation time is $T = 10$, and absolute errors on the slow variable have been computed with respect to the Euclidean norm of the difference in positions between analytical and numerical solutions. Stability is investigated using the same technique used in Subsection 5.1. Figures 3(a) and 3(b) illustrate errors as functions of mesostep $\delta$ and renormalized small step $\tau/\epsilon$. Observe that given $\delta$ errors are minimized at specific values of $\tau/\epsilon$ for both integrators, but the accuracy of nonintrusive FLAVOR is less sensitive to $\tau/\epsilon$. Figures 3(c) and 3(d) plot the optimal value of $\tau/\epsilon$ as a function of $\delta$ and the associated to error. Observe also that for nonintrusive FLAVOR the dependence of the optimal value of $\tau/\epsilon$ on $\delta$ is weak, whereas for artificial FLAVOR the optimal value of $\tau/\epsilon$ roughly scales linearly with $\delta$. Figure 3(e) and 3(f) describe how error changes with smallstep $\tau$ for mesostep $\delta$ fixed. Figure 3(e) can be viewed in correspondence with the condition $\delta << \tau/\epsilon$ required for accuracy. This requirement, however, is just a sufficient condition to obtain an error bound, as we can see in Figure 3(f). There the weak dependence of the error on $\tau/\epsilon$ for a fixed $\delta$ shows that one does not have to choose the microstep with too much care or optimize the integrator with respect to its value, if artificial FLAVOR is used. As a matter of fact, all the numerical experiments illustrated in this paper (except for Figures 3(c) and 3(d)) have been performed without any tuning of the value $\tau/\epsilon$. We have simply used the rule of thumb $\delta \sim \gamma \frac{\tau}{\epsilon}$ where $\gamma$ is a small parameter ($0.1$ for instance).

Therefore, it appears that the benefits of artificial FLAVORS lie in their superior accuracy and stability.

Notice that there is no resonant value of $\delta$ or $\tau$.

5.3 Numerical error analysis for nonlinear systems

In this subsection, we will consider the nonlinear Hamiltonian system

$$H(x, y, z, p_x, p_y, p_z) = \frac{1}{2}p_x^2 + \frac{1}{2}p_y^2 + \frac{1}{2}p_z^2 + x^4 + \epsilon^{-1} \frac{\omega_1}{2} (y - x)^2 + \epsilon^{-1} \frac{\omega_2}{2} (z - y)^2 \quad (5.8)$$

Thus, the potential is $U = \frac{\omega_1}{2} (y - x)^2 + \frac{\omega_2}{2} (z - y)^2$ and $V = x^4$. Here $\frac{x+y+z}{3}$ acts as a slow degree of freedom and $y - x$ and $z - y$ act as fast degrees of freedom.

Figure [5.8] illustrates $t \mapsto \frac{x(t) + y(t) + z(t)}{3}$ (slow variable, convergent strongly) and $t \mapsto (y(t) - x(t), z(t) - y(t))$ (fast variables, convergent in measure) computed with symplectic Euler and with the induced symplectic FLAVOR (2.5)). Define $q := (x, y, z)$. To
(a) Error of nonintrusive FLAVOR as a function of $\delta$ and $\tau/\sqrt{\epsilon}$. Notice that not all pairs of step lengths lead to stable integrations.

(b) Error of artificial FLAVOR as a function of $\delta$ and $\tau/\sqrt{\epsilon}$

(c) Optimal $\tau/\sqrt{\epsilon}$ and error of nonintrusive FLAVOR as functions of $\delta$

(d) Optimal $\tau/\sqrt{\epsilon}$ and error of artificial FLAVOR as functions of $\delta$

(e) Error dependence on $\tau/\sqrt{\epsilon}$ for a given $\delta$: nonintrusive FLAVOR

(f) Error dependence on $\tau/\sqrt{\epsilon}$ for a given $\delta$: artificial FLAVOR

Figure 3: Error analysis of (5.1). Parameters are $\omega = \sqrt{\epsilon} = 10^3$, $x(0) = 0.8$ and $y(0) = x(0) + 1.1/\omega$. 

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31
Figure 4: Comparison between trajectories integrated by Variational Euler and FLAVOR (defined by (2.5) and (2.9)). FLAVOR uses mesostep $\delta = 0.01$ and microstep $\tau = 0.0005$, and Symplectic Euler uses time step $\tau = 0.0005$. Time axes in the right column are zoomed in (by different ratios) to illustrate the fact that fast variables are captured in the sense of measure. FLAVOR accelerated the computation by roughly 20x ($\delta = 20\tau$). In this experiment $\epsilon = 10^{-6}$, $\omega_1 = 1.1$, $\omega_2 = 0.97$, $x(0) = 0.8$, $y(0) = 0.811$, $z(0) = 0.721$, $p_x(0) = 0$, $p_y(0) = 0$ and $p_z(0) = 0$. Simulation time $T = 50$.

Illustrate the $F$-convergence property of FLAVOR, we fix $H = 1$, vary the mesostep $\delta = H/M$ by changing $M$ and show the Euclidean norm error of the difference between $\frac{1}{M} \sum_{i=0}^{M-1} q(T - ih/M)$ computed with FLAVOR and computed with symplectic Euler in Figure 5(a). Notice that without an averaging over time length $h$, the error will be no longer monotonically but oscillatorily decreasing as $\delta$ changes (plots not shown), because fast variables are captured only in the sense of measure. As shown in Figure 5(a) the error scales linearly with $\frac{1}{M}$ for $M$ not too small, and therefore the global error is a linear function of the mesostep $\delta$ and the method is first order convergent. Figure 5(b) shows that the error in general grows linearly with the total simulation time, and this linear growth of the error has been observed for a simulation time larger than $\omega (\epsilon^{-1/2})$. Figure 5(c) shows that the error does not depend on $\omega (\epsilon^{-1/2})$ for a fixed $\delta$, as long as $\epsilon$ is not too large (i.e. $\omega$ not too small); the error is instead controlled by $M$. This is not caused by reaching the limit of machine accuracy, it is a characteristic of the method: the plateau for large $\omega$ corresponds to the complete scale separation regime of FLAVOR as a multiscale method.

Notice that there is no resonant value of $\delta$ in the sense of convergence.
The fact that the error scales linearly with total simulation time is a much stronger (numerical) result than our (theoretical) error analysis for FLAVORS (in which the error is bounded by a term growing exponentially with the total simulation time). We conjecture that the linear growth of the error is a consequence of the fact that FLAVOR is symplectic and is only true for a subclass of systems, possibly integrable systems. A rigorous analysis of the effects of the structure preservation of FLAVORS on long term behavior remains to be done.

6 Numerical experiments

6.1 Hidden Van der Pol oscillator (ODE)

Consider the following system ODEs

\[
\begin{align*}
\dot{r} &= \frac{1}{\epsilon} (r \cos \theta + r \sin \theta - \frac{1}{3} r^3 \cos^3 \theta) \cos \theta - \epsilon r \cos \theta \sin \theta \\
\dot{\theta} &= -\epsilon \cos^2 \theta - \frac{1}{\epsilon} (\cos \theta + \sin \theta - \frac{1}{3} r^2 \cos^3 \theta) \sin \theta
\end{align*}
\]  

(6.1)

where \( \epsilon \ll 1 \). Taking the transformation from polar coordinates to Cartesian coordinates by \([x, y] = [r \sin \theta, r \cos \theta]\) as the local diffeomorphism, we obtained the hidden system:

\[
\begin{align*}
\dot{x} &= -\epsilon y \\
\dot{y} &= \frac{1}{\epsilon} (x + y - \frac{1}{3} y^3)
\end{align*}
\]  

(6.2)

Taking the second time derivative of \( y \), the system can also be written as the 2nd-order ODE:

\[
\ddot{y} + y = \frac{1}{\epsilon} (1 - y^2) \dot{y}.
\]  

(6.3)

The latter is a classical Van der Pol oscillator [113]. Nonintrusive FLAVOR as defined by (1.34) can be directly applied to (6.1) (with hidden slow and fast processes) by turning on and off the stiff parameter \( \frac{1}{\epsilon} \). More precisely, defining \( \Phi^\epsilon_\alpha (r, \theta) \) by

\[
\Phi^\epsilon_\alpha (r, \theta) := \begin{pmatrix} r \\ \theta \end{pmatrix} + \alpha h \begin{pmatrix} r \cos \theta + r \sin \theta - \frac{1}{3} r^3 \cos^3 \theta \cos \theta \\ -(\cos \theta + \sin \theta - \frac{1}{3} r^2 \cos^3 \theta) \sin \theta \end{pmatrix} - \epsilon h \begin{pmatrix} r \cos \theta \sin \theta \\ \cos^2 \theta \end{pmatrix}
\]  

(6.4)
Figure 6: Over a timespan of $5/\epsilon$ (a) Direct Forward Euler simulation of (6.2) with time steps resolving the fast time scale (b) (nonintrusive (1.34)) FLAVOR simulation of (6.2) (c) Polar to cartesian image of the (nonintrusive (1.34)) FLAVOR simulation of (6.1) with hidden slow and fast variables. Forward Euler uses time step $h = 0.05\epsilon = 0.00005$. The two FLAVORS simulations use $\delta = 0.01$ and $\tau = 0.00005$. Parameters are $\frac{1}{\epsilon} = 1000$, $x(0) = 1$, $y(0) = 1$.

FLAVOR is defined by (1.34) with $\bar{u} := (\bar{r}, \bar{\theta})$, i.e.,

$$(\bar{r}_t, \bar{\theta}_t) = \left(\Phi_{\delta-\tau}^{0,\epsilon} \circ \Phi_{\tau}^{1,\epsilon}\right)^k(r_0, \theta_0) \quad \text{for} \quad k\delta \leq t < (k + 1)\delta. \quad (6.5)$$

We refer to Figure 6 for a comparison of integrations by Forward Euler, used as a benchmark, and FLAVORS. FLAVORS gives trajectories close to Forward Euler and correctly captures the $O(\frac{1}{\epsilon})$ period [13] of the relaxation oscillation. Moreover, a 200x acceleration is achieved using FLAVOR.

### 6.2 Hamiltonian system with nonlinear stiff and soft potentials

In this subsection, we will apply the Symplectic Euler FLAVOR defined by (2.5) and (2.9) to the mechanical system whose Hamiltonian is

$$H(y, x, p_y, p_x) := \frac{1}{2}p_y^2 + \frac{1}{2}p_x^2 + \epsilon^{-1}y^6 + (x - y)^4 \quad (6.6)$$

Here, stiff potential $\epsilon^{-1}U = \epsilon^{-1}y^6$ and soft potential $V = (x - y)^4$ are both nonlinear.
Figure 7: In this experiment, \( \epsilon = 10^{-6} \), \( y(0) = 1.1 \), \( x(0) = 2.2 \), \( p_y(0) = 0 \) and \( p_x(0) = 0 \). Simulation time \( T = 2 \). FLAVOR (defined by (2.5) and (2.9)) uses mesostep \( \delta = 10^{-3} \) and microstep \( \tau = 10^{-5} \), Variational Euler uses small time step \( \tau = 10^{-5} \), and IMEX uses mesostep \( \delta = 10^{-3} \). Since the fast potential is nonlinear, IMEX is an implicit method and nonlinear equations have to be solved at every step, and IMEX turns out to be slower than Variational Euler. FLAVOR is strongly accurate with respect to slow variables and accurate in the sense of measures with respect to fast variables. Comparing to Symplectic Euler, FLAVOR accelerated the computation by roughly 100x.

Figure 7 illustrates \( t \mapsto y(t) \) (dominated by a fast process), \( t \mapsto x(t) - y(t) \) (a slow process modulated by a fast process), and \( t \mapsto H(t) \) computed with: Symplectic Euler, the induced symplectic FLAVOR (2.5 and 2.9), and IMEX [105]. Notice that \( x - y \) is not a purely slow variable but contains some fast component, and therefore the FLAVOR integration of it contains a modulation of local oscillations, which could be interpreted as that fast component slowed down by FLAVOR. It’s not easy to find a purely slow variable or a purely fast variable in the form of (1.2) for this example, but the integrated trajectory for such a slow variable will not contain these slowed-down local oscillations.

Figure 8: Fermi-Pasta-Ulam problem [42] – 1D chain of alternatively connected harmonic stiff and non-harmonic soft springs
By Variational Euler with small time step $\tau' = 5 \times 10^{-5} = 0.05/\omega$. 38 periods in Subplot2 with zoomed-in time axis ($\sim 380$ in total over the whole simulation span).

By artificial FLAVOR (2.11) with mesostep $\delta = 0.002$ and microstep $\tau = 10^{-4} = 0.1/\omega$. 38 periods in Subplot2 axis ($\sim 380$ in total over the whole simulation span).

Figure 9: Simulations of the FPU problem over $T = 2\omega$. Subplot2 of both figures have zoomed-in time axes so that whether phase lag or any other distortion of trajectory exists could be closely investigated. In this experiment $m = 3$, $\omega = 10^3$, $x(0) = [0.4642, -0.4202, 0.0344, 0.1371, 0.0626, 0.0810]$ is randomly chosen and $y(0) = [0, 0, 0, 0, 0, 0]$.

6.3 Fermi-Pasta-Ulam problem

In this subsection, we will consider the Fermi-Pasta-Ulam (FPU) Problem [42] illustrated by Figure 8 and associated with the Hamiltonian

$$H(q, p) := \frac{1}{2} \sum_{i=1}^{m} (p_{2i-1}^2 + p_{2i}^2) + \frac{\omega^2}{4} \sum_{i=1}^{m} (q_{2i} - q_{2i-1})^2 + \sum_{i=0}^{m} (q_{2i+1} - q_{2i})^4. \quad (6.7)$$

The FPU problem is a well known benchmark problem [79, 54] for multiscale integrators because it exhibits different behaviors over widely separated timescales. The stiff springs nearly behave like harmonic oscillators with period $\sim O(\omega^{-1})$. Then, the centers of masses linked by stiff springs (i.e., the midpoints of stiff springs) change over a timescale $O(1)$. The third timescale, $O(\omega)$, is associated with the rate of energy exchange between stiff springs. The fourth timescale, $O(\omega^2)$, corresponds to the synchronization of energy exchange between stiff springs. On the other hand, the total energy of the stiff springs behaves almost like a constant. This wide separation of timescales can be seen in Figure 9 [10] and [12] where four subplots address different scales: Subplot1 shows the fast variables $(q_{2i} - q_{2i-1})/\sqrt{2}$; Subplot2 shows one of the slow variables $(q_{2i} + q_{i})/\sqrt{2}$; Subplot3 shows the energy transfer pattern among stiff springs, which is even slower; Subplot4 shows the near-constant total energy of three stiff springs. All four subplots are time-series. A comprehensive survey on FPU problem, including discussions on timescales and numerical recipes, can be found in [54].

Figures 9(a) and 9(b) compare symplectic Euler (with time steps fine enough to resolve FPU over the involved long time scale) and with the artificial FLAVOR (2.11). On a timescale $O(\omega)$ ($\omega \gg 1$), FLAVOR captured slow variable’s periodic behavior with
the correct period and phase, as well as the slower process of energy transfer. At the same time, FLAVOR accelerated the computation by roughly 40x (since $\delta = 40\tau'$).

It is not worrisome that artificial FLAVOR produces stiff spring energy trajectories with rapid local oscillations, which exhibit in both thicker individual energy curves and total energy with larger variance. In fact, these local oscillations do not seem to affect the global transfer pattern nor its period and are caused by the numerical error associated with microstep $\tau$. This can be inferred by using the artificial FLAVOR introduced in Subsection 2.1.4 with $\theta_\tau$ corresponding to the exact flow of $H^{fast}$ (rather than its Variational Euler approximation: this specific artificial Euler resembles the Impulse Method, but the Impulse Method will yield unbounded trajectories if one runs even longer time simulations, whereas FLAVORS do not seem to have an error growing exponentially with the total simulation time). As illustrated in Figure 11 exact flow helps to obtain thin energy curves of stiff springs with no rapid local oscillations as well as a total energy with a variance smaller than that given by fine Variational Euler (Figure 9(a)), with the transfer pattern similar to Figure 9(b).

Now, we reach further to $O(\omega^2)$ total integration time to investigate different integrators’ performances in capturing the synchronized energy exchange pattern (Figure 10).

There is a significant difference among stiff spring energy transfer patterns produced by Velocity Verlet, FLAVOR, IMEX and the Impulse Method. Here, there is no analytic solution or provably accurate method for comparison. FLAVOR is the only method that shows periodic behavior on the long time scale and convergence tests show that FLAVOR’s trajectories remain stable under small variations of step sizes. Notice that mathematically it can be shown that the dynamical system admits periodic orbits (for example, by Poincaré-Bendixson theorem). Furthermore, it is physically intuitive that the three stiff springs should alternatively obtain their maximal and minimal energies, and these maximal energies should be of fixed values. In addition, if we change the slow potential to be quadratic the system is still very similar to non-harmonic FPU but now analytically solvable. There, the energy exchanging pattern (Figure 12) resembles FLAVORS’ result of the non-harmonic system but not the other integrators’. Notice that if run on the modified quadratic FPU problem however, FLAVORS, Velocity-Verlet, IMEX and the Impulse Method all obtain perfect results (plots omitted). These are numerical evidences supporting FLAVOR on the $O(\omega^2)$ timescale.

It is worth discussing why Velocity-Verlet with a time step much smaller than the characteristic length of the fast scale ($O(1/\omega)$) is still not satisfactory. Being a second order method, it has an error bound of $O(e^T h^2)$. On the other hand, backward error analysis guarantees that the energy of the integrated trajectory oscillates around the true conserved energy, hence eliminating the possibility of exponential growth of the numerical solution. Nevertheless, at this moment there is no result known to the authors to link these two analytical results to guarantee long term accuracy on the stiff springs’ energies. The energy exchange among stiff springs is in fact an delicate phenomenon, and a slight distortion in stiff spring lengths could easily disrupt its period or even its periodicity.
These numerical observations seem to indicate that symplectic FLAVORS may have special long time properties. Specifically, although we could not quantify the error here because there is no benchmark to compare to when the total simulation time is $O(\omega^2)$, the long term behavior seems to indicate an error growing much slower than exponentially (please refer to Remark 1.7 for a discussion on exponential error bounds and Figure 5(b) for another example of conjectured linear error growth). A rigorous investigation on FLAVORS’ long time behavior remains to be done.

Figure 13 summarizes FLAVOR’s performance on various timescales in a comparison to Velocity Verlet.

Notice that there are many sophisticated methods designed for integrating the FPU problem (see [54] for a review), as well as general multiscale methods that can be applied to the FPU problem. HMM as one state-of-art method in the latter category, together
Figure 11: By artificial FLAVORS (Subsection 2.1.4) based on exact fast flow with mesostep \( \delta = 0.002 \) and microstep \( \tau = 10^{-4} \). Less oscillatory stiff spring energies. 38 periods in Subplot2 with zoomed-in time axis (∼380 in total over the whole simulation span).

with identification of slow variables [5] captured the energy transfer between stiff springs over a time span of the order of \( \omega \). Simulations shown here are over a time span of the order of \( O(\omega^2) \).

6.3.1 On resonances

Multiscale in time integrators are usually plagued by two kinds of resonances.

The first kind, called Takens resonance [107], is related to the fact there are no closed equations on slow variables [15]. FLAVORS avoid Takens resonance because, thanks to \( F \)-convergence, the information on the local invariant measure of fast variables is not lost. Observe that the FPU problem has Takens resonance (the eigenfrequencies of the strong potential are identical). Nevertheless, FLAVORS still capture the solution trajectories given any large value of \( \omega \) with mesostep \( \delta \gg 1/\omega \) independent of \( \omega \).

The second kind [25] is related to instabilities created by interactions between parameters \( \epsilon, \tau \) and \( \delta \). For instance, if \( \epsilon^{-1} = \omega^2 \) resonance might happen at \( \omega \delta \) or \( \omega \tau \) equal to multiples of \( \pi/2 \). The analysis provided in Section 5 shows that such unstable interaction does not occur, either in the sense of stability or in terms of peaks of error function. This can be intuitively understood by observing that FLAVORS never approximate \( \cos(\delta \omega) \), while on the other hand, it does approximate \( \cos(\tau \omega) \) whose resonance frequency \( \tau = 2\pi/\omega \) is ruled off by the requirement of \( \tau \ll \epsilon \) for nonintrusive FLAVOR and \( \tau \ll \sqrt{\epsilon} \) for artificial FLAVOR.

6.4 Nonlinear 2D primitive molecular dynamics

Now consider a two-dimensional, two degrees of freedom example in which a point mass is linked through a spring to a massless fixed hinge at the origin. While the spring as well
Figure 13: Quantities of interest in integrations of FPU over different timescales. FLAVOR \(2.11\) captures the fastest timescale in the sense of measure, while Velocity Verlet cannot accurately capture the slowest \(O(\omega^2)\) timescale despite of the small time step it uses. Here FLAVOR is 200 times faster than Velocity Verlet. All parameters are the same as in Figure 10(a) and 10(b), e.g. \(\omega = 200, \delta = 0.002, \tau = 0.0005\) and \(h = 10^{-5}\).

as the point mass are allowed to rotate around the hinge (the spring remains straight), the more the spring-mass tilts away from its equilibrium angle the more restorative force it will experience. This example is a simplified version of prevailing molecular dynamics models, in which bond lengths and angles between neighboring bonds are both spring-like; other potential energy terms are ignored.

Denote by \(x\) and \(y\) the Euclidean coordinates of the mass, and \(p_x, p_y\) corresponding momentums. Also introduce polar coordinates \((r, \theta)\), with \(x = r \cos \theta\) and \(y = r \sin \theta\). Then the Hamiltonian reads

\[
H = \frac{1}{2} p_x^2 + \frac{1}{2} p_y^2 + \frac{1}{2} \omega^2 (r - r_0)^2 + (\cos \theta)^2 \\
= \frac{1}{2} p_r^2 + \frac{1}{2} p_\theta^2 + \frac{1}{2} \omega^2 (\sqrt{x^2 + y^2} - r_0)^2 + \frac{x^2}{x^2 + y^2}
\]

where \(r_0\) is equilibrium bond length parameter and \(\omega\) is large number denoting bond oscillation frequency.

Remark 6.1. This seemingly trivial example is not easy to integrate.

1. If the system is viewed in Euclidean coordinates \((x, y, p_x, p_y)\) it is completely nonlinear with a nonpolynomial potential, and hence the Impulse Method or its variations [51, 110, 44, 96], or IMEX [105], or the homogenization method introduced in [20] cannot be applied using a mesostep.

2. If the Hamiltonian is rewritten in generalized coordinates \((r, \theta, p_r, p_\theta)\), \(H = \frac{1}{2} p_r^2 + \frac{1}{2} \omega^2 (r - r_0)^2 + \frac{1}{2} \cos(\theta)^2\), a fast quadratic potential can be identified.
However, the mass matrix $\begin{bmatrix} 1 & 0 \\ 0 & r^2 \end{bmatrix}$ is not constant, but rapidly oscillating, and hence methods that work for quasi-quadratic fast potentials (i.e. “harmonic oscillator” with a slowly changing frequency) ([20] for example) cannot be applied.

![Figure 14: Simulation of (6.8). Symplectic Euler uses small time step $\tau = 0.0002$ and the induced symplectic FLAVOR ((2.5) and (2.9)) uses mesostep $\delta = 0.01$ and microstep $\tau = 0.0002$. In this simulation $\omega = 500$, $x(0) = 1.1$, $y(0) = 0.8$, $p_x(0) = 0$, $p_y(0) = 0$ and simulation time $T = 100$. FLAVOR reproduced the slow $\theta$ trajectory while accelerating the simulation time by roughly 50x (since $\delta = 50\tau$). It can also be seen from both energy fluctuations and the trajectory of the fast variable that the fast processes’ amplitude is well captured although its period has been lengthened.]

6.5 Nonlinear 2D molecular clipper

We now consider a united-atom representation of a three atom polymer with two bonds (e.g. propane or water molecule). This is a simplified version of several prevailing molecular dynamics force fields (for example, CHARMM [21], AMBER [33], or a simpler example of butane [92, 93]). Using conservation of momentum, we fix the coordinate system in the 2D plane defined by the three atoms. Introduce both Cartesian coordinates $(x_1, y_1, x_2, y_2, x_3, y_3)$, as well as generalized coordinates $r_1 = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}$
and \( r_2 = \sqrt{(x_3 - x_2)^2 + (y_3 - y_2)^2} \) for bond lengths and \( \theta \) for the angle between the two bonds (Figure 15). The kinetic energy is

\[
K.E. = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{1}{2} m_2 (\dot{x}_2^2 + \dot{y}_2^2) + \frac{1}{2} m_3 (\dot{x}_3^2 + \dot{y}_3^2)
\]

(6.9)

where \( m_1, m_2, \) and \( m_3 \) denote the masses of the atoms.

The potential energy consists of a bond term and a bond angle term, both of which are of harmonic oscillator type:

\[
P.E. = V_{bond} + V_{angle}
\]

(6.10)

\[
V_{bond} = \frac{1}{2} K_r (r_1 - r_0)^2 + (r_2 - r_0)^2)
\]

(6.11)

\[
V_{angle} = \frac{1}{2} K_\theta (\cos(\theta) - \cos(\theta_0))^2
\]

(6.12)

Notice that the system is in fact fully nonlinear: if written in generalized coordinates, the kinetic energy will correspond to a nonlinear and position dependent mass matrix, whereas in Cartesian coordinates, both terms in the potential energy are non-polynomial functions in positions.

In the case of propane, \( m_1 = 15 \mu, m_2 = 14 \mu, m_3 = 15 \mu \) where \( \mu = 1.67 \cdot 10^{-27} kg, \)

\( r_0 = 1.53 \AA, K_r = 83.7 \text{kcal/(mol \AA}^2\), \( \theta_0 = 109.5^\circ \) and \( K_\theta = 43.1 \text{kcal/mol} \) [92].

The propane system is characterized by a separation of timescales to some extent: bond stretching and bond-angle bending are characterized by \( 10^{14} \) and \( 10^{13} \) Hz vibrational frequencies respectively [117]. For investigation on FLAVORS, we use unitless parameters and exaggerate the timescale separation by setting \( K_r \) to be 8370 and \( K_\theta \) to be 431. We also let \( \mu = 1 \) without loss of generality for arithmetic considerations.

In this system, the bond potential is the fast potential and the bond-angle potential is the slow one. It is well known that using a large time step at the timescale corresponding to the bond-angle potential by freezing bond lengths produces biased results, and many
physics based methods have been proposed to remedy this difficulty (for example by the
approach of Fixman \cite{Fixman}; also see a review in \cite{117}). On the other hand, few multiscale
methods work for this fully nonlinear system.

Figure 16: Simulations of exaggerated propane molecule (Subsection 6.5). Symplectic Eu-
er uses $h = 0.01$ and the induced symplectic FLAVOR ((2.5) and (2.9)) parameters are $\delta =
0.1$ and $\tau = 0.01$. Initial conditions are $[x_1, y_1, x_2, y_2, x_3, y_3] = [0, 1.533, 0, 2.6136, 1.0826]$ and
$[m_1, \dot{x}_1, m_1 \dot{y}_1, m_2, \dot{x}_2, m_2 \dot{y}_2, m_3, \dot{x}_3, m_3 \dot{y}_3] = [-0.4326, -1.6656, 0.1253, 0.2877, -1.1465, 1.1909]$.

Figure 16 compares symplectic Euler with the induced symplectic FLAVOR ((2.5) and (2.9)) applied in Euclidean coordinates. 10x acceleration is achieved. A simulation movie is also available at \url{http://www.cds.caltech.edu/~mtao/Propane.avi} and \url{http://www.acm.caltech.edu/~owhadi/}

6.6 Forced nonautonomous mechanical system: Kapitza’s inverted pen-
dulum

As the famous Kapitza’s inverted pendulum shows \cite{63} (for recent references see \cite{6} for
numerical integration and \cite{97} for generalization to the stochastic setting), the up
position of a single pendulum can be stabilized if the pivot of the pendulum experiences
external forcing in the form of vertical oscillation. Specifically, if the position of the
pivot is given by $y = \sin(\omega t)$, the system is governed by

$$
\dot{\theta} = [g + \omega^2 \sin(2\pi \omega t)] \sin \theta
$$

\hfill (6.13)
where $\theta$ denotes the clockwise angle of the pendulum from the positive $y$ direction, $l$ is the length of the pendulum and $g$ is the gravitational constant. In this case, the rapid vibration causes the pendulum to oscillate slowly around the positive $y$ direction with a $O(1)$ frequency.

Figure 17: Simulations of the inverted pendulum. The integration by Variational Euler + d’Alembert principle uses time step $h = 0.2/\omega/\sqrt{l} \approx 0.000067$, while FLAVOR (defined by (6.15)) uses $\delta = 0.002$ and $\tau = 0.2/\omega/\sqrt{l}$. Also, $g = 9.8$, $l = 9$, $\theta(0) = 0.2$, $\dot{\theta}(0) = 0$ and $\omega = 1000$

A single scale integration of this system could be done by Variational Euler with discrete d’Alembert principle for external forces

\[
\begin{align*}
  f_i &= \omega^2 \sin(2\pi \omega h) \\
  p_{i+1} &= p_i + h[g + f_i] \sin \theta_i \\
  \theta_{i+1} &= \theta_i + h p_{i+1}/l
\end{align*}
\]  

(6.14)

where the time step length $h$ has to be smaller than $O(1/\omega)$.

FLAVOR is given by

\[
\begin{align*}
  q_{n\delta+\tau} &= q_{n\delta} + \tau p_{n\delta}/l \\
  p_{n\delta+\tau} &= p_{n\delta} + \tau g \sin(q_{n\delta+\tau}) + \omega^2 \sin(2\pi \omega n \tau) \\
  q_{(n+1)\delta} &= q_{n\delta+\tau} + (\delta - \tau)p_{n\delta+\tau}/l \\
  p_{(n+1)\delta} &= p_{n\delta+\tau} + (\delta - \tau)g \sin(q_{(n+1)\delta})
\end{align*}
\]  

(6.15)
Observe that the time dependent force is synchronized on the \( \tau \) time scale instead of the \( \delta \) time scale, specifically \( \omega^2 \sin(2\pi \omega n \tau) \) instead of \( \omega^2 \sin(2\pi \omega n \delta) \) in (6.15).

Numerical results are illustrated in Figure 17 (also available as a movie at http://www.cds.caltech.edu/~mtao/InvertedPendulum.avi and http://www.acm.caltech.edu/~owhadi/). Notice in this example that \( \theta \), being the only degree of freedom, contains a combination of slow and fast dynamics. FLAVOR could only capture the fast dynamics in the sense of measures, and this is why dents appear as modulation on the slow oscillation of \( \theta \). On the other hand, although this forced system does not admit a conserved energy, the value of the Hamiltonian should oscillate periodically due to the periodic external driving force. While a non mechanics based method such as Forward Euler often produces an unbounded growth or a decrease in the energy, FLAVORS do not have this drawback.

**Remark 6.2.** Consider the case of a rapid potential of the form \( \Omega^2 (q_1) q_2^2 / \epsilon^2 \) (where \( q_1 \) is the slow an \( q_2 \) the fast variable). In the limit of a vanishing \( \epsilon \), it is known that the term contributes to the effective Hamiltonian with a contribution \( V(q_1) \) (the so called Fixman term). On may have the intuition that FLAVOR would only be consistent with a term of the form \( \gamma V(q_1) \), where \( 0 < \gamma < 1 \) is only a fraction of one, because the rapid force is only accounted for over a time \( \tau < \delta \). This intuition is not correct because the effect of FLAVOR is not to account the rapid force for over a time \( \tau < \delta \) but to slow down the rapid force by a fraction \( \tau/\epsilon \). This effect can also be seen in the algorithm (6.15) where the force term \( \omega^2 \sin(2\pi \omega n \tau) \) has been slowed down by a factor \( \tau/\delta \) (the Kapitza’s inverted pendulum illustrates a similar phenomenon where rapid oscillations contribute a stabilizing term to the effective Hamiltonian, nevertheless FLAVORS remain accurate).

### 6.7 Nonautonomous SDE system with hidden slow variables

Consider the following artificial nonautonomous SDE system

\[
\begin{aligned}
\frac{du}{dt} &= \frac{4}{3(u+v)^2} \left( -\frac{1}{2} \left( \frac{u-v}{2} \right)^2 + 5 \sin(2\pi t) \right) dt - \frac{1}{\sqrt{2}} \left( \left( \frac{u+v}{2} \right)^3 + c - \frac{u-v}{2} \right) dt - \sqrt{\frac{2}{\epsilon}} dW_t \\
\frac{dv}{dt} &= \frac{4}{3(u+v)^2} \left( -\frac{1}{2} \left( \frac{u-v}{2} \right)^2 + 5 \sin(2\pi t) \right) dt + \frac{1}{\sqrt{2}} \left( \left( \frac{u+v}{2} \right)^3 + c - \frac{u-v}{2} \right) dt + \sqrt{\frac{2}{\epsilon}} dW_t
\end{aligned}
\]  

(6.16)

where \( c \) is a positive constant and the two \( dW_t \) terms refer to the same Brownian Motion. The system (6.16) can be converted via the local diffeomorphism

\[
\begin{aligned}
u &= (x-c)^{1/3} - y \\
v &= (x-c)^{1/3} + y
\end{aligned}
\]

(6.17)

into the following hidden system separating slow and fast variables

\[
\begin{aligned}
dx &= -\frac{1}{2} y^2 dt + 5 \sin(2\pi t) dW_t \\
dy &= \frac{1}{3} (x-y) dt + \sqrt{\frac{2}{\tau}} dW_t
\end{aligned}
\]

(6.18)

Nonintrusive FLAVOR (3.7) can be directly applied to (6.16) using a time step \( \delta \gg \epsilon \) without prior identification of the slow and fast variables, i.e., without prior identification.
Figure 18: (a) Integration of (6.16) by noninvasive FLAVOR (3.7) using mesostep step $\delta = 0.01$ (b) Integration of (6.16) by Euler-Maruyama using fine time step $h = 10^{-4}$ (c) Integration of (6.18) by Euler-Maruyama using the same small step $h = 10^{-4}$. Expectations of the slow variable (whether or not hidden) are obtained by empirically averaging over an ensemble of 100 independent sample trajectories. $\epsilon = 10^{-4}$, $x(0) = 1 + \epsilon$, $y(0) = 1$, $T = 2$ (the expectation of the real solution will blow up around $T = 3$). We have chosen $c = 10$ so that the transformation is a diffeomorphism.

of the slow variable $x$ or of the system (6.18). The expected values of solutions of (6.16) integrated by FLAVORS with mesostep $\delta$ and Euler-Maruyama with a small time step $\tau$ are presented in Figure 18. FLAVOR has accelerated the computation by 100x.

6.8 Langevin equations with slow noise and friction

In this subsection, we consider the one dimensional, two degrees of freedom system modeled by the SDEs (now both springs are quartic rather than harmonic):

\[
\begin{aligned}
dy &= p_y dt \\
dx &= p_x dt \\
dp_y &= -\epsilon^{-1}y^3 dt - 4(y - x)^3 dt - cp_y dt + \sigma dW^1_t \\
dp_x &= -4(x - y)^3 dt - cp_x dt + \sigma dW^2_t.
\end{aligned}
\]  

(6.19)

We compare several autocorrelation functions and time-dependent moments of this stochastic process integrated by quasi-symplectic FLAVOR (4.7) and Geometric Langevin Algorithm (GLA) [17]. FLAVOR and GLA gave results in agreement
Figure 19: SDE (6.19): autocorrelation functions of $E[y(t)y(0)]$ (dominantly fast) and of $E[(x(t) - y(t))(x(0) - y(0))]$ (dominantly slow), empirically obtained by GLA and FLAVORS.

(Figure 19, 20(a), 20(b)). Since GLA is weakly-convergent and Boltzmann-Gibbs preserving, this is numerical evidence that quasi-symplectic FLAVOR is also.

Expectations are empirically calculated by averaging over an ensemble of 100 sample trajectories with $T = 30$, $\epsilon = 10^{-8}$, $\tau = 0.001$, $\delta = 0.01$. $y(0) = 2.1/\omega$ (with $\omega := 1/\sqrt{\epsilon}$), $x(0) = y(0) + 1.8$, $c = 0.1$ and $\sigma = 0.5$. GLA uses time step $h = 0.001$. Noise and friction are slow here in the sense that they are not of the order $O(\omega)$ or larger.

As shown in the plots, in the regime dominated by deterministic dynamics (roughly from $t = 0$ to $t = 8$) various moments calculated empirically by FLAVORS and GLA are in agreement, indicating that the same rate of convergence towards the Boltzmann-Gibbs distribution is obtained. And in that regime, autocorrelation functions of the slow variables agree, serving as numerical evidence that FLAVORS is weakly converging towards the SDE solution, whereas autocorrelation functions of the fast variables agree only in the sense of measures (after time averaging over a mesoscopic ($\ll(1)$) time span).

The fluctuations between FLAVORS and GLA for large time are an effect of the finite number of samples (100) used to compute sample averages.

Recall that if the noise is applied to slow variables, FLAVORS do not converge strongly but only in the sense of distributions.
Consider a system with the same configuration as above. The difference is that the soft spring oscillates at a frequency nonlinearly dependent on the stiff spring’s length, and the left mass experiences strong friction and noise while the right mass does not. The Hamiltonian is
\[ H(y, x, p_y, p_x) = \frac{1}{2} p_y^2 + \frac{1}{2} p_x^2/2 + \frac{1}{4} \omega^4 y^4 + e^y(x - y)^2, \] (6.20)
and the governing SDEs are:
\[
\begin{align*}
    dy &= p_y dt \\
    dx &= p_x dt \\
    dp_y &= -\omega^4 y^3 dt - (2 + y - x)(y - x)e^y dt - \omega^2 c p_y dt + \omega \sigma dB^t \\
    dp_x &= -2(x - y)e^y dt
\end{align*}
\] (6.21)

In this system, the deterministic dynamics and the effects of noise and friction both involve a $O(1/\omega^2)$ timescale. We have implemented the fast noise and friction version of FLAVORS ((4.8) and (2.9)).

In Figure 21, we have plotted the first and second moments of the dominantly slow variable $x(t) - y(t)$ as well as the first moment of the dominantly fast variable $y(t)$ as functions of time. Moments of the dominantly slow variable integrated by quasi-symplectic FLAVOR (Subsection 4.2.1) and GLA [17] concur, numerically suggesting weak convergence and preservation of Boltzmann-Gibbs. 100x computational acceleration is achieved.
Figure 21: $E|y(t) - y(t)|$, $E|y(t)|$, and $E|x(t) - y(t)|^2$ obtained by GLA and quasi-symplectic FLAVOR (Subsection 4.2.1). Expectations are empirically calculated by averaging over an ensemble of 50 sample trajectories with $T = 10$, $\omega = 100$, $\tau = 10^{-4}$, $\delta = 0.01$. $y(0) = 1/\omega$, $x(0) = y(0) + 1.8$, $c = 0.1$ and $\sigma = 1$. GLA uses time step $h = 10^{-4}$.

7 Appendix

7.1 Proof of theorems 1.1 and 1.2

Define the process $t \mapsto (\bar{x}_t, \bar{y}_t)$ by

$$ (\bar{x}_t, \bar{y}_t) := \eta(\bar{u}_t). \quad (7.1) $$

It follows from the regularity of $\eta$ that it is sufficient to prove the $F$-convergence of $(\bar{x}_t, \bar{y}_t)$ towards $\delta X_t \otimes \mu(X_t, dy)$. Moreover, it is also sufficient to prove the following inequalities (7.2), (7.3) in order to obtain inequalities (1.24) and (1.25)

$$ |x_t^\epsilon - \bar{x}_t| \leq Ce^Ct \psi_1(u_0, \epsilon, \delta, \tau) \quad (7.2) $$

and

$$ \left| \frac{1}{T} \int_t^{t+T} \varphi(\bar{x}_s, \bar{y}_s) ds - \int_{\mathbb{R}^p} \varphi(X_t, y) \mu(X_t, dy) \right| \leq \psi_2(u_0, \epsilon, \delta, \tau, T, t)(\|\varphi\|_{L^\infty} + \|\nabla \varphi\|_{L^\infty}) \quad (7.3) $$

Now define $\psi_\tau^\epsilon$ by

$$ \psi_\tau^\epsilon(x, y) := \eta \circ \theta_\tau^\epsilon \circ \eta^{-1}(x, y) \quad (7.4) $$
Define $\psi^\eta_h$ by
\[ \psi^\eta_h(x, y) := \eta \circ \theta_h^G \circ \eta^{-1}(x, y) \quad (7.5) \]

**Proposition 7.1.** The vector fields $f$ and $g$ associated with the system of Equations (1.2) are Lipschitz continuous. We also have
\[ (\bar{x}_t, \bar{y}_t) = (\psi^\epsilon_{\delta - \tau} \circ \psi^\epsilon_{\tau})^k(x_0, y_0) \quad \text{for} \quad k\delta \leq t < (k + 1)\delta. \quad (7.6) \]

Moreover, there exists $C > 0$ such that for $h \leq h_0$ and $\frac{\tau}{\epsilon} \leq \tau_0$ we have
\[ |\psi^\epsilon_{\tau}(x, y) - (x, y) - \tau(g(x, y), 0) - \frac{\tau}{\epsilon}(0, f(x, y))| \leq C\left(\frac{\tau}{\epsilon}\right)^2 \quad (7.7) \]
and
\[ |\psi^\eta_h(x, y) - (x, y) - h(g(x, y), 0)| \leq Ch^2. \quad (7.8) \]

Furthermore, given $x_0, y_0$, the trajectories of $(x^\epsilon_t, y^\epsilon_t)$ and $(\bar{x}_t, \bar{y}_t)$ are uniformly bounded in $\epsilon, \delta \leq h_0, \tau \leq \min(\tau_0, \epsilon)$.

**Proof.** Since $(x, y) = \eta(u)$, we have
\[ \dot{x} = (G + \frac{1}{\epsilon}F)\nabla \eta^x \circ \eta^{-1}(x, y) \quad (7.9) \]
\[ \dot{y} = (G + \frac{1}{\epsilon}F)\nabla \eta^y \circ \eta^{-1}(x, y). \quad (7.10) \]

Hence, we deduce from Equation (1.2) of Condition 1.1 that
\[ g(x, y) = G\nabla \eta^x \circ \eta^{-1}(x, y) \quad (7.11) \]
\[ f(x, y) = F\nabla \eta^y \circ \eta^{-1}(x, y). \quad (7.12) \]

We deduce the regularity of $f$ and $g$ from the regularity of $G, F$ and $\eta$. Equation (7.6) is a direct consequence of the definition of $\psi^\epsilon_{\tau}$ and $\psi^\eta_h$ and Equation (1.27) (we write $(x_0, y_0) := \eta(u_0)$). Observe that Equation (1.2) of Condition 1.1 also requires that
\[ F\nabla \eta^x = 0 \quad G\nabla \eta^y = 0. \quad (7.13) \]

Now observe that
\[ \psi^\epsilon_{\tau}(x, y) - (x, y) - (g(x, y), 0)\tau - (0, f(x, y))\frac{\tau}{\epsilon} = (\eta \circ \theta^G_{\tau} - \eta \circ (G\nabla \eta^x, 0) - \frac{\tau}{\epsilon}(0, F\nabla \eta^y)) \circ \eta^{-1}(x, y). \quad (7.14) \]

Using (7.13), (1.29), Taylor expansion and the regularity of $\eta$ we obtain (7.7). Similarly
\[ \psi^\eta_h(x, y) - (x, y) - h(g(x, y), 0) := \left(\eta \circ \theta^G_h - \eta(x, y) - h\left(G\nabla \eta^x, 0\right)\right) \circ \eta^{-1}(x, y). \quad (7.15) \]

Using (7.13), (1.28), Taylor expansion and the regularity of $\eta$ we obtain (7.8). The uniform bound (depending on $x_0, y_0$) on the trajectories of $(x^\epsilon_t, y^\epsilon_t)$ and $(\bar{x}_t, \bar{y}_t)$ is a consequence of the uniform bound (given $u_0$) on the trajectories of $u^\epsilon_t$ and $\bar{u}_t$. \[ \square \]
It follows from Proposition 7.1 that it is sufficient to prove theorems 1.1 and 1.2 in the situation where \( \eta \) is the identity diffeomorphism. More precisely, the \( F \)-convergence of \( \bar{u}_t \) is a consequence of the \( F \)-convergence of \( (\bar{x}_t, \bar{y}_t) \) and the regularity of \( \eta \). Furthermore, from the uniform bound (depending on \( (x_0, y_0) \)) on the trajectories of \( (x_t, \bar{y}_t) \) and \( (\bar{x}_t, \bar{y}_t) \) we deduce that \( g \) and \( f \) are uniformly bounded and Lipschitz continuous (in \( \epsilon, \delta \leq h_0, \tau \leq \min(\eta_0, \delta) \)) over those trajectories.

Define
\[
\bar{g} := \int g(x, y) \mu(dx, dy)
\]
where \( \mu \) is the family of measures introduced in Condition 1.2. Let us prove the following lemma.

**Lemma 7.1.**

\[
|x^e_{n\delta} - \bar{x}_{n\delta}| \leq Ce^{Cn}\left( \delta + \left( \frac{\tau}{\epsilon} \right)^2 + \sup_{1 \leq l \leq n} |J(l)| \right) \tag{7.16}
\]

with \( J(k) = J_1(k) + J_2(k) \),

\[
J_1(k) := \sum_{n=0}^{k-1} \int_{n\delta}^{(n+1)\delta} g(x^e_{n\delta}, y^e_s) \, ds - \delta \bar{g}(x^e_{n\delta}) \tag{7.17}
\]

and

\[
J_2(k) := \sum_{n=0}^{k-1} \delta \left( \bar{g}(x_{n\delta}) - g(x_{n\delta}, y_{n\delta}) \right) \tag{7.18}
\]

**Proof.** Observe that

\[
x^e_{(n+1)\delta} = x^e_{n\delta} + \int_{n\delta}^{(n+1)\delta} g(x^e_{n\delta}, y^e_s) \, ds + \int_{n\delta}^{(n+1)\delta} (g(x^e_s, y^e_s) - g(x^e_{n\delta}, y^e_s)) \, ds \tag{7.19}
\]

Hence,

\[
x^e_{(n+1)\delta} - \bar{x}_{(n+1)\delta} = x^e_{n\delta} - \bar{x}_{n\delta} + I_1 + I_2(n) + I_3 + I_4(n) + I_5 \tag{7.20}
\]

with

\[
I_1 := \int_{n\delta}^{(n+1)\delta} (g(x^e_s, y^e_s) - g(x^e_{n\delta}, y^e_s)) \, ds \tag{7.21}
\]

\[
I_2(n) := \int_{n\delta}^{(n+1)\delta} g(x^e_{n\delta}, y^e_s) \, ds - \delta \bar{g}(x^e_{n\delta}) \tag{7.22}
\]

\[
I_3 := \delta \left( \bar{g}(x^e_{n\delta}) - \bar{g}(\bar{x}_{n\delta}) \right) \tag{7.23}
\]

\[
I_4(n) := \delta \left( \bar{g}(x_{n\delta}) - g(x_{n\delta}, y_{n\delta}) \right) \tag{7.24}
\]

\[
I_5 := \delta \bar{g}(x_{n\delta}, y_{n\delta}) - (\bar{x}_{(n+1)\delta} - \bar{x}_{n\delta}) \tag{7.25}
\]

51
Now observe that
\[ |I_1| \leq \| \nabla_x g \|_{L^\infty} \delta^2 \] (7.26)
and
\[ |I_3| \leq \delta \| \nabla_x g \|_{L^\infty} |x_{n\delta}^e - \bar{x}_{n\delta}|. \] (7.27)

Using (7.7) and (7.8) we obtain that
\[ |I_5| \leq C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) \] (7.28)

Combining the previous equations, we have obtained that
\[ x_{(n+1)\delta}^e - \bar{x}_{(n+1)\delta} \leq x_{n\delta}^e - \bar{x}_{n\delta} + C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) + C \delta |x_{n\delta}^e - \bar{x}_{n\delta}| + (I_2 + I_4)(n) \] (7.29)
and
\[ x_{(n+1)\delta}^e - \bar{x}_{(n+1)\delta} \geq x_{n\delta}^e - \bar{x}_{n\delta} - C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) - C \delta |x_{n\delta}^e - \bar{x}_{n\delta}| + (I_2 + I_4)(n) \] (7.30)

Write
\[ J(n) := \sum_{k=0}^{n-1} (I_2 + I_4)(k) \] (7.31)

Summing the first \( n \) inequalities (7.29) and (7.30)
\[ x_{n\delta}^e - \bar{x}_{n\delta} \leq C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) n + C \delta \sum_{k=0}^{n-1} |x_{k\delta}^e - \bar{x}_{k\delta}| + J(n) \] (7.32)
\[ x_{n\delta}^e - \bar{x}_{n\delta} \geq -C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) n - C \delta \sum_{k=0}^{n-1} |x_{k\delta}^e - \bar{x}_{k\delta}| + J(n) \] (7.33)

Hence
\[ |x_{n\delta}^e - \bar{x}_{n\delta}| \leq C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) n + C \delta \sum_{k=0}^{n-1} |x_{k\delta}^e - \bar{x}_{k\delta}| + |J(n)| \] (7.34)

And we obtain by induction
\[ |x_{n\delta}^e - \bar{x}_{n\delta}| \leq C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^2 \right) \left( n + C \delta \sum_{k=1}^{n} (n-k)(1+C\delta)^{k-1} \right) \] (7.35)
\[ |J(n)| + C \delta \sum_{l=2}^{n} (1+C\delta)^{l-2} |J(n-l+1)| \]

Equation (7.35) concludes the proof of Lemma 7.1.

We now need to control \( J_1(k) \) and \( J_2(k) \). First, let us prove the following lemma.
**Lemma 7.2.** For $N \in \mathbb{N}^*$ we have

$$|J_1(k)| \leq (\delta k)C\left(\frac{\delta e^{C\frac{\epsilon}{N\epsilon}}} + E\left(\frac{\delta}{N\epsilon}\right)\right) \quad (7.36)$$

*Proof.* Define $\hat{y}_t^\epsilon$ such that $\hat{y}_t^\epsilon = y_t^\epsilon$ for $t = (n + j/N)\delta$, $j \in \mathbb{N}^*$, and

$$\frac{d\hat{y}_t^\epsilon}{dt} = \frac{1}{\epsilon} f(x_{n\delta}^\epsilon, \hat{y}_t^\epsilon) \quad \text{for} \quad (n + j/N)\delta \leq t < (n + (j + 1)/N)\delta. \quad (7.37)$$

Using the regularity of $f$ and $g$ we obtain that

$$|\hat{y}_t^\epsilon - y_t^\epsilon| \leq C\delta e^{C\frac{\epsilon}{N\epsilon}}. \quad (7.38)$$

First, observe that

$$1 \leq \frac{1}{\delta} \int_{n\delta}^{(n+1)\delta} g(x_{n\delta}^\epsilon, y_s^\epsilon) \, ds - \bar{g}(x_{n\delta}^\epsilon) = K_1 + K_2 \quad (7.39)$$

with

$$K_1 := \frac{1}{\delta} \sum_{j=0}^{N-1} \int_{(n+j/N)\delta}^{(n+(j+1)/N)\delta} (g(x_{n\delta}^\epsilon, y_s^\epsilon) - g(x_{n\delta}^\epsilon, \hat{y}_s^\epsilon)) \, ds \quad (7.40)$$

and

$$K_2 := \frac{1}{N} \sum_{j=0}^{N-1} \left(\frac{N}{\delta} \int_{(n+j/N)\delta}^{(n+(j+1)/N)\delta} g(x_{n\delta}^\epsilon, \hat{y}_s^\epsilon) \, ds - \bar{g}(x_{n\delta}^\epsilon)\right). \quad (7.41)$$

We have

$$|K_1| \leq \|\nabla g\|_{L^\infty} \frac{1}{N} \sum_{j=0}^{N-1} \sup_{(n+j/N)\delta \leq s \leq (n+(j+1)/N)\delta} |y_s^\epsilon - \hat{y}_s^\epsilon|. \quad (7.42)$$

Hence, we obtain from (7.38) that

$$|K_1| \leq C\delta e^{C\frac{\epsilon}{N\epsilon}} \quad (7.43)$$

Moreover, we obtain from conditions 1.2 and 1.3 that

$$|K_2| \leq CE\left(\frac{\delta}{N\epsilon}\right) \quad (7.44)$$

This concludes the proof of Lemma 7.2. □

**Lemma 7.3.** We have for $m \in \mathbb{N}^*$

$$|J_2(k)| \leq C\delta k\left( m\delta + E\left(\frac{mr}{\epsilon}\right) + \left(\frac{r}{\epsilon} + m\delta + m\left(\frac{r}{\epsilon}\right)^2\right)e^{C\frac{m\epsilon}{\epsilon}}\right), \quad (7.45)$$

53
Proof. Let \( m \in \mathbb{N}^* \). Define \((\bar{x}_s, \bar{y}_s)\) such that for \( j \in \mathbb{N}^* \), \( n \in \mathbb{N}^* \),

\[
\begin{align*}
\frac{d\bar{x}_s}{dt} &= g(\bar{x}_s, \bar{y}_s) \quad \text{for} \quad jm\delta \leq s < (j+1)m\delta, \\
\frac{d\bar{y}_s}{dt} &= \frac{1}{\epsilon} f(\bar{x}_s, \bar{y}_s) \quad \text{for} \quad n\delta \leq s < n\delta + \tau, \\
\bar{y}_s &= \bar{y}_{n\delta + \tau} \quad \text{for} \quad n\delta + \tau \leq s < (n+1)\delta, \\
\bar{y}_{(n+1)\delta} &= \bar{y}_{n\delta + \tau} \quad \text{for} \quad n + 1 \neq jm, \\
(\bar{x}_{jm}, \bar{y}_{jm}) &= (\bar{x}_{jm\delta}, \bar{y}_{jm\delta}).
\end{align*}
\] (7.46)

Define \( \tilde{y}^a_s \) by

\[
\begin{align*}
\frac{d\tilde{y}^a_t}{dt} &= \frac{1}{\epsilon} f(\bar{x}_{jm\delta}, \tilde{y}^a_t) \quad \text{for} \quad jm\tau \leq t < (j+1)m\tau, \\
\tilde{y}^a_{jm\tau} &= \bar{y}_{jm\delta},
\end{align*}
\] (7.47)

and define \( \bar{x}^a_n \) by

\[
\bar{x}^a_n = \bar{x}_{jm\delta} \quad \text{for} \quad jm \leq n < (j+1)m.
\] (7.48)

Observe that

\[
J_2(k) = K_3 + K_4 + K_5 + K_6 + K_7
\] (7.49)

with

\[
K_3 := \sum_{n=0}^{k-1} \left( \int_{n\delta}^{(n+1)\delta} g(\bar{x}_s, \bar{y}_s) \, ds - \delta g(\bar{x}_{n\delta}, \bar{y}_{n\delta}) \right),
\] (7.50)

\[
K_4 := \sum_{n=0}^{k-1} \delta \left( \frac{1}{\tau} \int_{n\tau}^{(n+1)\tau} g(\bar{x}^a_n, \tilde{y}^a_s) \, ds - \int_{n\delta}^{(n+1)\delta} g(\bar{x}_s, \tilde{y}_s) \, ds \right),
\] (7.51)

\[
K_5 := \frac{\delta}{\tau} \sum_{n=0}^{k-1} \left( \tau \bar{g}(\bar{x}^a_n) - \int_{n\tau}^{(n+1)\tau} g(\bar{x}^a_n, \tilde{y}^a_s) \, ds \right),
\] (7.52)

\[
K_6 := \delta \sum_{n=0}^{k-1} \left( \bar{g}(\bar{x}_{n\delta}) - \bar{g}(\bar{x}^a_n) \right).
\] (7.53)

Using the regularity of \( g \) we obtain

\[
|K_6| \leq \delta kC\delta m.
\] (7.54)

Arranging the right hand side of (7.51) into groups of \( m \) terms corresponding to the intervals of (7.47) we obtain, from Condition 1.2 and Condition 1.3, that

\[
|K_5| \leq Ck\delta E\left( \frac{m\tau}{\epsilon} \right).
\] (7.55)

Using (7.48) and the regularity of \( f \) and \( g \) we obtain the following inequality

\[
|\tilde{y}^a_{jm\tau} - \tilde{y}_t| \leq Cm\delta e^{C\frac{m\tau}{\epsilon}}.
\] (7.56)
It follows that

\[ |K_4| \leq C\delta k m \delta e^{C \frac{m \tau}{\epsilon}}. \]  

(7.57)

Similarly, using (7.7) and (7.8), we obtain the following inequalities

\[ |\tilde{y}_{n\delta} - \bar{y}_{n\delta}| \leq C \left( \frac{\tau}{\epsilon} + m \delta + m \left( \frac{\tau}{\epsilon} \right)^2 \right) e^{C \frac{m \tau}{\epsilon}}, \]  

(7.58)

\[ |\tilde{x}_{n\delta} - \bar{x}_{n\delta}| \leq C m \left( \delta + \left( \frac{\tau}{\epsilon} \right)^2 \right). \]  

(7.59)

It follows that

\[ |K_3| \leq C \delta k \left( \frac{\tau}{\epsilon} + m \delta + m \left( \frac{\tau}{\epsilon} \right)^2 \right) e^{C \frac{m \tau}{\epsilon}}. \]  

(7.60)

This concludes the proof of Lemma 7.3.

Combining Lemma 7.1, 7.2 and 7.3 we have obtained that

\[ |x_{\epsilon n\delta} - \bar{x}_{n\delta}| \leq C e^{C \delta n} \left( \sqrt{\delta} + \left( \frac{\tau}{\epsilon} \right)^2 \frac{1}{\delta} + \delta e^{C \frac{\delta}{N \epsilon}} + E \left( \frac{\delta}{N \epsilon} \right) + E \left( \frac{m \tau}{\epsilon} \right) \right) \]  

(7.61)

Choosing \( N \) such that \( e^{C \frac{\delta}{N \epsilon}} \sim \delta^{-\frac{1}{2}} \) (observe that we need \( \epsilon \leq \delta/(-C \ln \delta) \)) and \( m \) such that \( \frac{m \tau}{\epsilon} e^{C \frac{m \tau}{\epsilon}} \sim \left( \frac{\delta}{\tau} + \frac{\tau}{\epsilon} \right)^{-\frac{1}{2}} \) we obtain for \( \frac{\delta}{\tau} + \frac{\tau}{\epsilon} \leq 1 \) that

\[ |x_{\epsilon n\delta} - \bar{x}_{n\delta}| \leq C e^{C \delta n} \left( \sqrt{\delta} + \left( \frac{\tau}{\epsilon} \right)^2 \frac{1}{\delta} + E \left( \frac{1}{C} \ln \frac{1}{\delta} \right) \right) \]  

(7.62)

This concludes the proof of inequality (7.2). The proof of (7.3) is similar and is also a consequence of (7.2).

7.2 Proof of Theorem 3.1

Define the process \( t \mapsto (\bar{x}_t, \bar{y}_t) \) by

\[ (\bar{x}_t, \bar{y}_t) := \eta(\bar{u}_t). \]  

(7.63)

It follows from the regularity of \( \eta \) that it is sufficient to prove the \( F \)-convergence of \((\bar{x}_t, \bar{y}_t)\) towards \( \delta X_t \otimes \mu(X_t, dy) \). Now define \( \psi_x^\epsilon \) by

\[ \psi_x^\epsilon(x, y, \omega) := \eta \circ \theta^\epsilon_x(\cdot, \omega) \circ \eta^{-1}(x, y), \]  

(7.64)

Define \( \psi_y^\tilde{p} \) by

\[ \psi_y^\tilde{p}(x, y, \omega) := \eta \circ \theta^{\tilde{p}}_y(\cdot, \omega) \circ \eta^{-1}(x, y). \]  

(7.65)
Proposition 7.2. The vector fields $f$, $g$ and matrix fields $\sigma$, $Q$ associated with the system of Equations (3.2) are uniformly bounded and Lipschitz continuous. We also have

$$\begin{cases}
(\bar{x}_0, \bar{y}_0) = \eta(u_0) \\
(\bar{x}(k+1)\delta, \bar{y}(k+1)\delta) = \psi^0_{\delta-k}((\bar{x}_{k\delta}, \bar{y}_{k\delta}), \omega_k) \\
(\bar{x}_t, \bar{y}_t) = (\bar{x}_{k\delta}, \bar{y}_{k\delta}) \quad \text{for} \quad k\delta \leq t < (k+1)\delta
\end{cases} \quad (7.66)$$

where $\omega_k, \omega'_k$ are i.i.d. samples from the probability space $(\Omega, F, \mathbb{P})$. Moreover there exists $C > 0$ and and $d$-dimensional centered Gaussian vectors $\xi'(\omega), \xi''(\omega)$ with identity covariance matrices such that for $h \leq h_0$ and $\tau \leq \tau_0$ we have

$$\left(\mathbb{E}\left[\left|\psi^0_{\delta-k}(x, y, \omega) - (x, y) - h(g(x, y), 0) - \sqrt{h}\sigma(x, y)\xi'(\omega), 0\right|^{2}\right]\right)^{\frac{1}{2}} \leq Ch^{\frac{3}{2}}, \quad (7.67)$$

$$\left(\mathbb{E}\left[\left|\psi^0_{\delta-k}(x, y, \omega) - (x, y) - \tau(g(x, y), 0) - \frac{\tau}{\epsilon}(0, f(x, y)) - \sqrt{\tau}\sigma(x, y)\xi''(\omega), 0\right|^{2}\right]\right)^{\frac{1}{2}} \leq C\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}}. \quad (7.68)$$

Proof. Since $(x, y) = \eta(u)$, we obtain from (3.1) and Itô’s formula

$$dx = \left((G + \frac{1}{\epsilon}F)\nabla\eta^x \circ \eta^{-1}(x, y)\right) dt + \left(\nabla\eta^x (H + \frac{1}{\sqrt{\epsilon}}K)\right) \circ \eta^{-1}(x, y) dW_t \quad + \frac{1}{2} \sum_{ij} \partial_i \partial_j \eta^x((H + \frac{1}{\sqrt{\epsilon}}K)(H + \frac{1}{\sqrt{\epsilon}}K)^T)_{ij} dt \quad (7.69)$$

$$dy = \left((G + \frac{1}{\epsilon}F)\nabla\eta^y \circ \eta^{-1}(x, y)\right) dt + \left(\nabla\eta^y (H + \frac{1}{\sqrt{\epsilon}}K)\right) \circ \eta^{-1}(x, y) dW_t \quad + \left(\frac{1}{2} \sum_{ij} \partial_i \partial_j \eta^y((H + \frac{1}{\sqrt{\epsilon}}K)(H + \frac{1}{\sqrt{\epsilon}}K)^T)_{ij}\right) \circ \eta^{-1} dt. \quad (7.70)$$

Hence we deduce from Equation (3.2) of Condition 3.1 that

$$g(x, y) = \left(G\nabla\eta^x + \frac{1}{2} \sum_{ij} \partial_i \partial_j \eta^x(HH^T)_{ij}\right) \circ \eta^{-1}(x, y) \quad (7.71)$$

$$\sigma(x, y) = (\nabla\eta^x H) \circ \eta^{-1}(x, y) \quad (7.72)$$

$$f(x, y) = \left(F\nabla\eta^y + \frac{1}{2} \sum_{ij} \partial_i \partial_j \eta^y(KK^T)_{ij}\right) \circ \eta^{-1}(x, y) \quad (7.73)$$

$$Q(x, y) = (\nabla\eta^y K) \circ \eta^{-1}(x, y). \quad (7.74)$$
Remark 7.1. Observe that Equation (3.2) of Condition 3.1 requires that
\[ F \nabla \eta^x = 0 \quad G \nabla \eta^y, \]  
(7.75)
\[ \sum_{ij} \partial_i \partial_j \eta^x (KK^T)_{ij} = 0, \]  
(7.76)
\[ \sum_{ij} \partial_i \partial_j \eta^y (HH^T)_{ij} = 0, \]  
(7.77)
\[ \sum_{ij} \partial_i \partial_j \eta^x (KH^T + HK^T)_{ij} = 0, \]  
(7.78)
and
\[ \sum_{ij} \partial_i \partial_j \eta^y (KH^T + HK^T)_{ij} = 0. \]  
(7.79)
Equations (7.78) and (7.79) are satisfied if \( KH^T \) is skew-symmetric. One particular case could be, of course, \( KH^T = 0 \), which translates into the fact that for all \( u \) the ranges of \( H(u) \) and \( K(u) \) are orthogonal, i.e., the noise with amplitude \( 1/\sqrt{\epsilon} \) is applied to degrees of freedom orthogonal to those with \( O(1) \) noise.

We deduce the regularity of \( f, g, \sigma \) and \( Q \) from the regularity of \( G, F, H, K \) and \( \eta \). Equation (7.6) is a direct consequence of the definition of \( \psi^\epsilon \) and \( \psi^h \) and Equation (7.66). Now observe that
\[ \psi^\epsilon(x, y, \omega) - (x, y) - \tau(g(x, y), 0) - \frac{\tau}{\epsilon}(0, f(x, y)) - \sqrt{\tau}(\sigma(x, y)\xi'(\omega), 0) \]
\[ - \sqrt{\tau}(0, Q(x, y)\xi'(\omega)) = \left( \eta \circ \theta^\epsilon_x - \eta - \tau(G \nabla \eta^x + \frac{1}{2} \sum_{ij} \partial_i \partial_j \eta^x (HH^T)_{ij}, 0) \right) \]
\[ - \frac{\tau}{\epsilon}(0, F \nabla \eta^y + \frac{1}{2} \sum_{ij} \partial_i \partial_j \eta^y (KH^T)_{ij}) - \sqrt{\tau}(\nabla \eta^y H\xi'(\omega), 0) \]
\[ - \sqrt{\tau}(0, \nabla \eta^y K\xi'(\omega)) \circ \eta^{-1}(x, y). \]  
(7.80)
Using Equations (7.75), (7.76), (7.77), (7.78) and (7.79), the Taylor-Ito expansion of \( \eta \circ \theta^\epsilon_x \), the regularity of \( \eta \), and Setting \( \xi'(\omega) \) equal to \( \xi \) defined in Equation (3.14) we obtain Equation (7.68). The proof of Equation (7.67) is similar.

It follows from Proposition 7.2 that it is sufficient to prove Theorem 3.1 in the situation where \( \eta \) is the identity diffeomorphism. More precisely the \( F \)-convergence of \( \bar{u}_t \) is a consequence of the \( F \)-convergence of \( (\bar{x}_t, \bar{y}_t) \) and the regularity of \( \eta \).

Let \( x \mapsto \varphi(x) \) be a function with continuous and bounded derivatives up to order 3. Let us prove the following lemma.
Lemma 7.4. We have
\[
E\left[\varphi(\bar{x}_{(n+1)\delta})\right] - E\left[\varphi(\bar{x}_{n\delta})\right] = \delta E\left[g(\bar{x}_{n\delta}, \bar{y}_{n\delta}) \nabla \varphi(\bar{x}_{n\delta}) + \sigma \sigma^T(\bar{x}_{n\delta}, \bar{y}_{n\delta}) : \text{Hess} \varphi(\bar{x}_{n\delta})\right] + I_0
\]  
(7.81)
with
\[
|I_0| \leq C\left(\delta^2 + \left(\frac{\tau}{\epsilon}\right)^2\right).
\]
(7.82)

Proof. Write \((\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau}) = \psi(\bar{x}_{n\delta}, \bar{y}_{n\delta}, \omega_n)\). Using Equation (7.68) we obtain that there exists an \(N(0, 1)\) random vector \(\xi_n\) independent from \((\bar{x}_{n\delta}, \bar{y}_{n\delta})\) and such that
\[
\bar{x}_{n\delta+\tau} - \bar{x}_{n\delta} = g(\bar{x}_{n\delta})\tau + \sqrt{\tau}\sigma(\bar{x}_{n\delta}, \bar{y}_{n\delta})\xi_n + I_1
\]
(7.83)
with
\[
(E[|I_1|^2])^{\frac{1}{2}} \leq C\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}}.
\]
(7.84)
Hence
\[
\left| E\left[\varphi(\bar{x}_{n\delta+\tau})\right] - E\left[\varphi(\bar{x}_{n\delta})\right] - \tau E\left[g(\bar{x}_{n\delta}, \bar{y}_{n\delta}) \nabla \varphi(\bar{x}_{n\delta})\right] + \sigma \sigma^T(\bar{x}_{n\delta}, \bar{y}_{n\delta}) : \text{Hess} \varphi(\bar{x}_{n\delta}) \right| \leq C\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}}
\]
(7.85)
Similarly, using Equation (7.67) we obtain that there exists an \(N(0, 1)\) random vector \(\xi'_n\) independent from \((\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau})\), and such that
\[
\bar{x}_{(n+1)\delta} - \bar{x}_{n\delta+\tau} = g(\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau})(\delta - \tau) + \sigma(\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau})\sqrt{\delta - \tau}\xi'_n + I_2
\]
(7.86)
with
\[
(E[|I_2|^2])^{\frac{1}{2}} \leq C(\delta - \tau)^{\frac{3}{2}}.
\]
(7.87)
Whence
\[
\left| E\left[\varphi(\bar{x}_{(n+1)\delta})\right] - E\left[\varphi(\bar{x}_{n\delta+\tau})\right] - (\delta - \tau)E\left[g(\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau}) \nabla \varphi(\bar{x}_{n\delta+\tau})\right] + \sigma \sigma^T(\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau}) : \text{Hess} \varphi(\bar{x}_{n\delta+\tau}) \right| \leq C(\delta - \tau)^{\frac{3}{2}}.
\]
(7.88)
Using the regularity of \(\sigma\), we obtain that
\[
\left( E\left[|\sigma(\bar{x}_{n\delta+\tau}, \bar{y}_{n\delta+\tau}) - \sigma(\bar{x}_{n\delta}, \bar{y}_{n\delta})|^2\right]\right)^{\frac{1}{2}} \leq C\left(\delta^4 + \sqrt{\frac{\tau}{\epsilon}}\right).
\]
(7.89)
The proof of (7.81) follows from (7.68), (7.85), (7.88), (7.89) and the regularity of \(g\) and \(\varphi\). □
Lemma 7.5. We have

\[
\left| \frac{E[\varphi(x_{n\delta})] - \varphi(x_0)}{n\delta} - L\varphi(x_0) \right| \leq J_5
\]  

(7.90)

with (for $\delta \leq C\tau/\epsilon$)

\[
|J_5| \leq C \left( \left( \frac{\delta\epsilon}{\tau} \right)^{\frac{1}{4}} + \left( \frac{\tau}{\epsilon} \right)^{\frac{1}{2}} + \frac{\sqrt{\tau}}{\epsilon} \right) + CE \left( \frac{1}{C} \ln \frac{\tau}{\delta\epsilon} \right).
\]  

(7.91)

Proof. Define $\tilde{B}_t$ by $\tilde{B}_0 = 0$ and

\[
\tilde{B}_t - \tilde{B}_{nt} = B_{n\delta + t} - B_{n\delta} \quad \text{for} \quad n\tau \leq t \leq (n+1)\tau.
\]  

(7.92)

Define $\tilde{y}_s$ by $\tilde{y}_0 = y_0$ and

\[
d\tilde{y}_t = \frac{1}{\epsilon} f(x_0, \tilde{y}_t) dt + \frac{1}{\sqrt{\epsilon}} Q(x_0, \tilde{y}_t) dB_t.
\]  

(7.93)

Write

\[
\bar{g}(x_0) := \int g(x_0, y) \mu(x_0, dy).
\]  

(7.94)

Using Lemma 7.4 we obtain

\[
E[\varphi(x_{n\delta})] - \varphi(x_0) = L\varphi(x_0) + J_1 + J_2 + J_3 + J_4,
\]  

(7.95)

with

\[
L\varphi(x_0) := \bar{g}(x_0) \nabla \varphi(x_0) + \bar{\sigma}\bar{\sigma}^T(x_0) : \text{Hess} \varphi(x_0),
\]  

(7.96)

\[
J_1 = \frac{1}{n} \sum_{k=0}^{n-1} E\left[ g(\bar{x}_{k\delta}, \bar{y}_{k\delta}) \nabla \varphi(\bar{x}_{k\delta}) + \sigma\sigma^T(\bar{x}_{k\delta}, \bar{y}_{k\delta}) : \text{Hess} \varphi(\bar{x}_{k\delta}) \right]
\]  

(7.97)

\[
J_2 = \frac{1}{n} \sum_{k=0}^{n-1} \left( E\left[ g(\bar{x}_0, \bar{y}_{k\delta}) \nabla \varphi(\bar{x}_0) + \sigma\sigma^T(\bar{x}_0, \bar{y}_{k\delta}) : \text{Hess} \varphi(\bar{x}_0) \right] \right.
\]  

\[- \frac{1}{\tau} \int_{k\tau}^{(k+1)\tau} E\left[ g(x_0, \bar{y}_s) \nabla \varphi(x_0) + \sigma\sigma^T(x_0, \bar{y}_s) : \text{Hess} \varphi(x_0) \right] ds \bigg),
\]  

(7.98)

\[
J_3 = \frac{1}{n\tau} \int_0^{n\tau} E\left[ g(x_0, \bar{y}_s) \nabla \varphi(x_0) + \sigma\sigma^T(x_0, \bar{y}_s) : \text{Hess} \varphi(x_0) \right] ds - L\varphi(x_0),
\]  

(7.99)

\[
|J_4| \leq C \left( \delta^2 + \left( \frac{\tau}{\epsilon} \right)^{\frac{1}{2}} \frac{1}{\delta} \right).
\]  

(7.100)
Using the regularity of $\sigma, g, \varphi$, (7.6) and (7.7) we obtain

$$|J_1| \leq C\left( (n\delta)^{\frac{3}{2}} + n\delta + n\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}} \right).$$  \hfill (7.101)

Using Property 3 of Condition 3.1 and Property 3 of Condition 3.2 we obtain

$$|J_3| \leq CE\left(\frac{n\tau}{\epsilon}\right).$$  \hfill (7.102)

Using (7.67) and (7.68), we obtain the following inequality

$$\left( \mathbb{E}\left[ |\bar{y}_{n\delta} - \bar{y}_{n\tau}|^2 \right] \right)^{\frac{1}{2}} \leq C\left( \sqrt{\frac{\tau}{\epsilon}} + (n\delta)^{\frac{3}{2}} + n\delta + n\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}} \right) \frac{n\tau}{\epsilon} e^{C\frac{n\tau}{\epsilon}},$$  \hfill (7.103)

which leads to

$$|J_2| \leq C\left( \sqrt{\frac{\tau}{\epsilon}} + (n\delta)^{\frac{3}{2}} + n\delta + n\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}} \right) e^{C\frac{n\tau}{\epsilon}}.$$  \hfill (7.104)

Hence, we have obtained

$$\left| \mathbb{E}\left[ \varphi(\bar{x}_{n\delta}) \right] - \varphi(x_0) \mathbb{E}\left[ \phi(\bar{x}_{n\delta}) \right] - L\varphi(x_0) \right| \leq J_5,$$  \hfill (7.105)

with

$$|J_5| \leq C\left( \sqrt{\frac{\tau}{\epsilon}} + (n\delta)^{\frac{3}{2}} + n\delta + n\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}} \right) e^{C\frac{n\tau}{\epsilon}} + E\left( \frac{n\tau}{\epsilon} \right) + C\left(\frac{\tau}{\epsilon}\right)^{\frac{3}{2}} \frac{1}{\delta}.$$  \hfill (7.106)

Choosing $n$ such that $\sqrt{\frac{n\tau}{\epsilon}} e^{C\frac{n\tau}{\epsilon}} \sim \left(\frac{\tau}{\epsilon}\right)^{\frac{1}{4}}$ we obtain (7.91) for $\delta \leq C\tau/\epsilon$. \qed

We now combine Lemma 7.5 with Theorem 1 of Chapter 2 of [104] which states that the uniform convergence (in $x_0$, $y_0$) of $\mathbb{E}\left[ \varphi(\bar{x}_{n\delta}) \right] - \varphi(x_0)$ to $L\varphi(x_0)$ as $\epsilon \downarrow 0, \tau \leq \delta, \frac{\tau}{\epsilon} \downarrow 0, \frac{\delta}{\tau} \downarrow 0$ and $(\frac{\tau}{\epsilon})^{\frac{3}{2}} \frac{1}{\delta} \downarrow 0$ implies the convergence in distribution of $\bar{x}_{n\delta}$ to the Markov process generated by $L$.

The $F$-convergence of $(\bar{x}_t, \bar{y}_t)$ can be deduced from the convergence in distribution of $\bar{x}_t$ and Equation (3.4) of Condition 3.1 The proof follows the same lines as above, which will not be repeated here.

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