Search for vertical stratification of metals in atmospheres of blue horizontal-branch stars

V. R. Khalack1, F. LeBlanc1, B. B. Behr2,**, G. A. Wade3, and D. Bohlender4

1 Département de Physique et d’Astronomie, Université de Moncton, Moncton, N.-B., E1A 3E9, Canada
e-mail: khalack@umoncton.ca
2 Department of Astronomy, University of Texas at Austin, 1 University Station C1400, Austin TX 78712-0259, USA
3 Department of Physics, Royal Military College of Canada, PO Box 17000 stn “FORCES”, Kingston, Ontario, K7K 4B4, Canada
4 National Research Council of Canada, Herzberg Institute of Astrophysics, 5071 West Saanich Road, Victoria, BC, V9E 2E7, Canada

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ABSTRACT

Context. The observed abundance peculiarities of many chemical species relative to the expected cluster metallicity in blue horizontal-branch (BHB) stars presumably appear as a result of atomic diffusion in the photosphere. The slow rotation (typically $v \sin i < 10 \text{ km s}^{-1}$) of BHB stars with effective temperatures $T_{\text{eff}} > 11500 \text{ K}$ supports this idea since the diffusion mechanism is only effective in a stable stellar atmosphere.

Aims. In this work we search for observational evidence of vertical chemical stratification in the atmospheres of six hot BHB stars: B84, B267 and B279 in M 15 and WF2-2541, WF4-3085 and WF4-3485 in M 13.

Methods. We undertake an abundance stratification analysis of the stellar atmospheres of the aforementioned stars, based on acquired Keck HIRES spectra.

Results. We have found from our numerical simulations that three stars (B267, B279 and WF2-2541) show clear signatures of the vertical stratification of iron whose abundance increases toward the lower atmosphere, while the other two stars (B84 and WF4-3485) do not. For WF4-3085 the stratification results are inconclusive. B267 also shows a signature of titanium stratification. Our estimates for radial velocity, $v \sin i$ and overall iron, titanium and phosphorus abundances agree with previously published data for these stars after taking the measurement errors into account. The results support the hypothesis regarding the efficiency of atomic diffusion in the stellar atmospheres of BHB stars with $T_{\text{eff}} > 11500 \text{ K}$.

Key words. stars: atmospheres – stars: horizontal-branch – stars: chemically peculiar

1. Introduction

According to the current understanding of stellar evolution, the horizontal-branch (HB) stars are post-main sequence stars that burn helium in their core and hydrogen in a shell (e.g. Moehler 2004). In this paper we consider the HB stars that are located in the blue part of the HB, to the left of the RR Lyrae instability strip. Most researchers call them blue horizontal-branch (BHB) stars to distinguish them from the red horizontal-branch (RHB) stars, which exhibit different observational properties. Sandage & Wellerstein (1960) have found from analysis of the colour–magnitude diagrams of globular clusters that the HB generally becomes bluer with decreasing metallicity. Derived masses of the cool ($T_{\text{eff}} < 11500 \text{ K}$) BHB stars in the globular cluster NGC 6388 (Moehler & Sweigart 2006) are in a good agreement with the predictions of canonical HB evolution, except for the hot BHB stars with $T_{\text{eff}} > 11500 \text{ K}$, where the estimated stellar masses seem to be lower than the canonical values.

The Hertzsprung–Russell diagrams of some globular clusters show long blue tails (an extension of the HB), populated by very hot BHB stars and extreme horizontal branch (EHB) stars. Published data on the BHB stars argue that the hot BHB stars show remarkable differences in physical properties when compared to the cool BHB stars. Using high-precision photometry of stars in M 13, Ferraro et al. (1998) have found gaps in the distribution of stars along the blue tail. One of these gaps, labeled as G1, is located at $T_{\text{eff}} \sim 11000–12000 \text{ K}$. Grundahl et al. (1998) used the results of Strömgren $u$b$γ$β$–photometry finding good agreement between the theoretical prediction of stellar evolution models and the observed location of BHB stars, except for the hot BHB stars, whose $u$-magnitudes are brighter than predicted. It appears that this $u$-jump is observed for the hot BHB stars and coincides with the temperature range of the G1 gap in M 13. Similar $u$-jumps have also been found for other globular clusters (Grundahl et al. 1999). For the hot BHB with effective temperatures up to 20000 K, the surface gravities derived from the fits of Balmer and helium line profiles appear to be lower than the predictions of stellar evolution models, while the gravities derived for the stars outside this temperature range are in good agreement with theoretical predictions (Moehler et al. 1995, 1997a,b, 2003). The stellar rotation velocity distribution of BHB stars also appears to have a discontinuity at $T_{\text{eff}} \approx 11500 \text{ K}$ (Peterson et al. 1995; Behr et al. 2000a; Recio-Blanco et al. 2004), indicating that the hotter stars show modest rotation with $v \sin i < 10 \text{ km s}^{-1}$, while the cooler stars rotate more rapidly.
on average. Comprehensive surveys of abundances also show that the hot BHB stars have abundance anomalies when compared to the cool BHB stars in the same cluster (Glaspey et al. 1989; Grundahl et al. 1999; Behr et al. 1999, 2000b, Behr 2003a; Fabbian et al. 2005; Pace et al. 2006).

The observed phenomena such as the low gravity, photometric jumps and gaps, abundance anomalies and slow rotation suggest that atomic diffusion could be important in the stellar atmospheres of hot BHB stars. Atomic diffusion arises from the competition between radiative acceleration and gravitational settling. This can produce a net acceleration on atoms and ions, which results in their diffusion through the atmosphere (Michaud 1970). In order for atomic diffusion to produce a vertical stratification of the abundances of particular elements, the stellar atmosphere must be hydrodynamically stable. According to Landstreet (1998), photospheric convection should be very weak at the effective temperatures of BHB stars. Theoretical atmospheric models of Hui-Bon-Hoa et al. (2000) showed that the observed photometric jumps and gaps for hot BHB stars can be explained by elemental diffusion in their atmosphere. Behr (2003b) has shown that adoption of a microturbulent velocity of 0 or 1 km s\(^{-1}\) provides the best fit to line strengths in the spectra of hot BHB stars. This fact supports the proposal that strong velocity fields are not present in the atmospheres of hot BHB stars.

While synthesizing spectral line profiles, Khalack et al. (2007) have recently found vertical abundance stratification of sulfur in the atmosphere of the field BHB star HD 135485. In this paper we also attempt to detect signatures of vertical abundance stratification of elements from line profile analyses of several other BHB stars for which we have appropriate data. Together with the data on stratification of the sulfur abundance in HD 135485, new positive results would provide a convincing argument in favour of efficient atomic diffusion in the atmospheres of hot BHB stars. In Sect. 2 we discuss the properties of the acquired spectra, while in Sect. 3 we describe details concerning the simulation routine and adopted atmospheric parameters for the program stars. The evidence for vertical stratification of some chemical species is given in Sect. 4, while the estimation of mean abundances and velocities is described in Sect. 5. A discussion follows in Sect. 6.

2. Observations

In this paper we have selected hot BHB stars from the list of objects published by Behr (2003b) and which have comparatively high signal-to-noise ratio (S/N) spectra available. Spectroscopic observations of the selected stars were undertaken in August 1998 with the Keck I telescope and the HIRES spectrograph. The journal of spectroscopic observations is shown in Table 1 where individual columns give the object identification, the heliocentric Julian Date of the observation, the exposure time, the S/N per pixel, the spectral coverage, the size of the seeing disk (FWHM) and the size of the C1 slit that provides a spectral resolution of R = \(\lambda/\delta\lambda\) = 45 000. For the aforementioned stars, Behr (2003b) has found that the underfilling of the slit should not change the estimated spectral resolution by more than 4–7%.

The package of routines developed by McCarthy (1990) for the FIGARO data analysis package (Shortridge 1993) was employed to process the spectra. A comprehensive description of the data acquisition and reduction procedure is presented by Behr (2003b). This followed a standard prescription with bias subtraction, flat-fielding, order extraction, and wavelength calibration from thorium-argon comparison lamp observations. To minimize the potential distortion of narrow spectra features, cosmic ray hits were identified and removed by hand.

### Table 1. Journal of Keck+HIRES spectroscopic observations of the selected hot BHB stars from Behr (2003b).

| Cluster/Star | HJD 245 000+ | Exposure time (s) | S/N | Coverage (\(\AA\)) | Seeing (arcsec) | Slit width (arcsec) |
|--------------|-------------|-----------------|-----|-------------------|----------------|-------------------|
| M13/WF2-2541 | 1046.7902   | 3 x 1500        | 44  | 3885–6292        | 0.90           | 0.86              |
| M13/WF4-3085 | 1052.7338   | 3 x 1200        | 37  | 3888–5356        | 1.10           | 0.86              |
| M13/WF4-3485 | 1053.7793   | 3 x 1200        | 34  | 3888–5356        | 0.90           | 0.86              |
| M15/B84      | 1053.9024   | 4 x 1400        | 34  | 3888–5356        | 0.80           | 0.86              |
| M15/B267     | 1053.9731   | 4 x 1400        | 26  | 3888–5356        | 0.80           | 0.86              |
| M15/B279     | 1053.8312   | 4 x 1400        | 34  | 3888–5356        | 0.90           | 0.86              |

3. Line profile simulations

3.1. Stellar atmosphere parameters

The line profile simulations were performed using the ZEEMAN2 spectrum synthesis code (Landstreet 1988; Wade et al. 2001). The stellar atmosphere models were calculated with the PHOENIX code (Hauschildt et al. 1997) assuming LTE (Local Thermodynamic Equilibrium) and using the stellar atmosphere parameters extracted from Behr (2003b) and listed in Table 2. In calculating the stellar atmosphere models for program stars, we have used solar metallicity with the enhanced iron and depleted helium abundances derived by Behr (2003b). The depleted helium abundance has also been taken into account during line profile simulations using ZEEMAN2. We used Gaussian instrumental profiles with widths derived from the comparison arc spectra.

To simulate the spectra of BHB stars, Behr (2003b) initially adopted a microturbulent velocity \(\xi = 2\text{ km s}^{-1}\) and subsequently updated it during the simulation routine according to the results of his model fits. He has obtained best fit values of \(\xi\) from 0 to 0.9 km s\(^{-1}\) for the BHB stars studied here (see Table 2). Since we search for possible vertical abundance stratification, our results can be affected by microturbulence. To estimate this influence we have adopted here two different microturbulent velocities \(\xi = 0\) and 2 km s\(^{-1}\) for our simulations.

### Table 2. Parameters of stellar atmospheres for the selected hot BHB stars from Behr (2003b).

| Cluster/Star | \(T_{\text{eff}}\) (K) | \(\log g\) (dex) | \(\xi\) (km s\(^{-1}\)) | \(V_1\) (km s\(^{-1}\)) |
|--------------|----------------|----------------|----------------|----------------|
| M13/WF2-2541 | 13 000        | 4.0            | 0.0            | -257.5         |
| M13/WF4-3085 | 14 000        | 4.0            | 0.0            | -255.7         |
| M13/WF4-3485 | 13 000        | 4.0            | 0.5            | -246.8         |
| M15/B84      | 12 000        | 3.5            | 0.5            | -108.2         |
| M15/B267     | 11 000        | 3.5            | 0.0            | -114.4         |
| M15/B279     | 11 000        | 3.5            | 0.9            | -104.4         |
Table 3. List of iron spectral lines used for the abundance analysis.

| \( \lambda, \, \text{Å} \) | \( \log gf \) | \( E_i, \, \text{cm}^{-1} \) | \( \log y_{\text{wd}} \) | Ref. | Object name |
|-----------------|----------------|-----------------|-----------------|---------|-------------|
| WF2-2541        | WF4-3085       | WF2-3585        | B84             | B267    | B279        |
| 4122.668        | \(-3.300\)     | 20830.58        | 8.49            | R&U     | ...         |
| 4128.748        | \(-3.578\)     | 20830.58        | 8.61            | R&U     | ...         |
| 4173.461        | \(-2.617\)     | 20830.58        | 8.61            | R&U     | ...         |
| 4177.618        | \(-3.776\)     | 73395.93        | 8.61            | R&U     | ...         |
| 4177.692        | \(-3.449\)     | 20516.96        | 8.61            | R&U     | ...         |
| 4178.862        | \(-2.535\)     | 20830.58        | 8.49            | R&U     | ...         |
| 4184.261        | \(-1.938\)     | 90397.87        | 8.35            | R&U     | ...         |

Table 3 is presented in full in electronic form. A portion is shown here for guidance and content.

### 3.2. Procedure

We examined the spectrum of each star to establish a line list suitable for stratification and stratiﬁcation analysis. The line identiﬁcation was performed using the VALD-2 (Kupka et al. 1999; Ryabchikova et al. 1999), Pickering et al. (2001) and Raassen & Uylings (1998) line databases. For all of the program stars the best represented element (with a number of readily visible line proﬁles) is iron. Some stars also show Ti ii or P ii lines. Therefore iron, titanium and phosphorus were selected for analysis of their possible vertical abundance stratification. For Fe ii lines, we extracted atomic data from Raassen & Uylings (1998), while for P ii lines we used VALD-2. Atomic data for Ti ii lines were taken from Pickering et al. (2001).

The ZEEMAN2 code has been modiﬁed (Khalack & Wade 2006) to allow for an automatic minimization of the model parameters using the downhill simplex method (Press et al. 1992). The minimization routine ﬁnds the global minimum of the \( \chi^2 \) function, which is speciﬁed as the measure of differences between simulated and observed line proﬁles. The relatively poor efﬁciency of the downhill simplex method, requiring a large number of function evaluations, is a well known problem. However, repeating the minimization routine several times in the vicinity of a supposed minimum in the chosen parameter space allows us to verify if the method converges to the global minimum (for more details about this minimization routine see Khalack et al. 2007).

To search for the presence of abundance stratiﬁcation, we have estimated the abundance of a chemical element from an analysis of their possible vertical abundance stratiﬁcation. The detected stratiﬁcation of iron, obtained in some of the stars studied here, is far too large (>2 dex) to be interpreted as the result of measurement errors. Our stellar atmosphere models were calculated taking into account the enhanced iron abundance and helium depletion (Behr 2003b) and the error in estimating the iron abundances, for a given model atmosphere, is expected to be less than 0.2 dex. Meanwhile, Khan & Shulyak (2007) have shown that using atmospheric models with varying Fe abundance (from one to ten times solar) can modify the abundances inferred from observed line proﬁles by up to ±0.25 dex. From the results in Table 4 we can conclude that no detectable iron stratification exists in WF4-3485 and B84.

### 4. Vertical abundance stratiﬁcation

Applying the technique described in Sect. 3.2, we have attempted to determine if the iron abundance is vertically stratiﬁed in the atmospheres of the six selected BHB stars. In comparison, we have also investigated the possible vertical stratiﬁcation of titanium and phosphorus for some BHB stars, where these elements are represented by a sufﬁcient number of spectral lines.

All of the selected BHB stars are slowly rotating and have sharp absorption lines that result in only 8 to 13 spectral bins per line proﬁle. This fact, together with the S/N and the uncertainties in the atomic data, are the main contributors to errors in the abundance inferred from a single line. This makes the detection of weak vertical abundance variations difﬁcult.

Our abundance analysis for \( \xi = 0 \, \text{km s}^{-1} \) shows that in the atmospheres of four BHB stars (B267, B279, WF4-3085 and WF2-2541) the iron abundance generally increases towards the deeper atmospheric depths (see Figs. 1–5). It should be noted that the range of optical depths diagnosed by the iron lines is much smaller in the stars B279 and WF2-2541 than in the other two stars shown here. To estimate the stratiﬁcation proﬁle of the iron abundance we ﬁt the data with a straight line using a least-square algorithm. We have also checked the dependence of each element’s abundance with respect to the lower level excitation potential \( E_i \) for the sample of analyzed lines (see as an example Fig. 3). Slopes of the linearly approximated dependence of each element’s abundance with respect to \( \log E_i \) and \( \log T_{5000} \) are given in Table 4 for the case of zero microturbulent velocity. Table 4 shows a clear correlation between the slope of the abundance versus \( \log T_{5000} \) and \( \log E_i \) for the elements considered in each star. This correlation strongly suggests that the obtained stratiﬁcation is physical and not due to uncertainties in the description of the temperature proﬁle of the stellar atmosphere.

The detected stratiﬁcation of iron, obtained in some of the stars studied here, is far too large (>2 dex) to be interpreted as the result of measurement errors. Our stellar atmosphere models were calculated taking into account the enhanced iron abundance and helium depletion (Behr 2003b) and the error in estimating the iron abundances, for a given model atmosphere, is expected to be less than 0.2 dex. Meanwhile, Khan & Shulyak (2007) have shown that using atmospheric models with varying Fe abundance (from one to ten times solar) can modify the abundances inferred from observed line proﬁles by up to ±0.25 dex.

From the results in Table 4 we can conclude that no detectable iron stratification exists in WF4-3485 and B84.
stars (B267 and B279) show quite large slopes of iron abundance with respect to both $\log E_\xi$ and $\log \tau_{3000}$. The corresponding slopes for the stars WF2-2541 and WF2-3085 are smaller but still statistically significant. The least significant of these slopes is that of the abundance versus $\log E_\xi$ for WF2-3085 which has only approximately a $3\sigma$ value.

Another factor that must be evaluated is the effective temperature of the underlying atmospheric model used. To verify the potential importance of this factor, we calculated the Fe stratification with models with an effective temperature higher and lower by 1000 K from the $T_{\text{eff}}$ values listed in Table 2. We also calculated the stratification with models assuming solar He and Fe abundances to evaluate the effect of this change on the inferred stratification profiles. Our simulations show that models with higher $T_{\text{eff}}$ and models with solar abundances usually slightly decrease the slope of the abundance with respect to $\log E_\xi$, while models with lower $T_{\text{eff}}$ increase this slope for all elements. In the case of iron in WF4-3085, the model with higher $T_{\text{eff}}$ and the model with solar abundances result in a negligible slope of its abundance versus $\log E_\xi$. Therefore, the results are not sufficient to report unambiguous detection of vertical stratification of iron in WF4-3085. The slopes for iron found for the other three stars (B267, B279 and WF-2541) are still significant in all these models, and stratification is therefore confirmed.

The two BHB stars (B84 in M 15 and WF4-3485 in M 13) which do not show any signs of stratification of their iron abundances are not different from the other stars with regards to their stellar atmosphere characteristics and rotation (see Table 2). However, these two stars possess very small titanium and phosphorus abundances as compared to the other stars studied here. Their spectra have very weak or invisible TiII and PII lines. Their iron abundance is also close to the solar value, at least for $\xi = 0 \text{ km s}^{-1}$ (see Table 5).

We observe an upturn in the iron abundance obtained assuming $\xi = 0 \text{ km s}^{-1}$ at low optical depths for some of the stars studied (e.g. Figs. 1a, 2a, 4a, 5a). The points on the figures corresponding to the upturn feature are the strong FeII lines and are not taken into account during the least-square fit of the data. To estimate the influence of microturbulence we have performed an abundance analysis with $\xi = 2 \text{ km s}^{-1}$, which leads to elimination of the upturn (see, for example, Fig. 1a). Performing a set of simulations with different microturbulent velocities, we have determined the minimum value of the microturbulence ($\xi_{\text{min}}$) for which the upturn disappears for each star studied. These values are reported in Table 5. As the inclusion of microturbulent str
velocity amplifies the iron stratification, the stratification obtained for iron in the first four stars studied with $\xi = 0$ km s$^{-1}$ can be considered to be a lower limit to the possible stratification in these stars. A more detailed study including abundance stratification in the calculation of the synthesized spectra might shed some light on this strange behaviour, but this is outside of the scope of the present paper.

Figure 1b indicates that B267 shows a possible signature of vertical stratification of the titanium abundance. The respective slopes of Ti abundance versus $\log E_i$ and $\log \tau_{5000}$ are significantly higher than zero for this star and are statistically significant (see Table 4). For B279, we can not confidently conclude that stratification exists since the slope of the abundance with respect to $\log E_i$ becomes weak when using either solar abundances or assuming a $T_{\text{eff}}$ increase of 1000 K. The slope of Ti abundance with respect to $\log \tau_{5000}$ also becomes weak in the model assuming a $T_{\text{eff}}$ increase of 1000 K. It should be noted that the Ti lines in B279 do not sample a large portion of the atmosphere and this fact renders the detection of stratification more difficult.

**Table 4.** Slopes of vertical abundance stratification calculated for the linear least-square fit to the data assuming zero microturbulence.

| Cluster/Star | Slopes of $\log (\text{Fe}/N_{\text{tot}})$ vs. $\log E_i$ | $n$ | Slopes of $\log (\text{P}/N_{\text{tot}})$ vs. $\log E_i$ | $n$ |
|--------------|--------------------------------------------------------|-----|--------------------------------------------------------|-----|
| M 13/WF2-2541 | 0.18 $\pm$ 0.04, 0.17 $\pm$ 0.03 | 53  | -0.67 $\pm$ 0.31, -0.14 $\pm$ 0.06 | 16  |
| M 13/WF4-3085 | 0.10 $\pm$ 0.03, 0.15 $\pm$ 0.02 | 117 | 0.54 $\pm$ 0.31, 0.24 $\pm$ 0.04 | 12  |
| M 13/WF4-3485 | -0.06 $\pm$ 0.03, -0.02 $\pm$ 0.03 | 50  | ... | ... |
| M 15/B84     | 0.04 $\pm$ 0.05, 0.03 $\pm$ 0.04 | 25  | ... | ... |
| M 15/B267    | 0.49 $\pm$ 0.06, 0.33 $\pm$ 0.03 | 46  | 0.47 $\pm$ 0.10, 0.36 $\pm$ 0.07 | 16  |
| M 15/B279    | 0.23 $\pm$ 0.04, 0.23 $\pm$ 0.02 | 71  | 0.21 $\pm$ 0.11, 0.30 $\pm$ 0.08 | 11  |
It is clear that phosphorous shows no clear signs of stratification in WF2-2541. For WF4-3085 we cannot make firm conclusions concerning phosphorous stratification since its abundance variation in the range of optical depths under consideration is only approximately 0.4 dex. Also, the slope of its abundance with respect to log $E_i$ is not statistically significant.

5. Mean abundances and velocities

The mean photospheric abundances derived for the selected BHB stars from an analysis of the available iron, titanium and phosphorus lines are reported in Table 5. For each chemical element, three columns are shown containing the mean abundance for simulations with microturbulent velocity $\xi = 0$ and 2 km s$^{-1}$ respectively, and the number of analyzed line profiles. The reported uncertainties are equal to the standard deviation calculated from the results of individual line simulations for all lines considered.

The derived heliocentric radial and projected rotation velocities as well as microturbulent velocities (see Table 5) are generally in agreement with the data provided by Behr (2003b) for these stars. Some inconsistencies are found between $\varv \sin i$ values obtained for WF2-2541 ($\varv \sin i = 0.0^{+0.07}_{-0.04}$ km s$^{-1}$) and B279 ($\varv \sin i = 5.92^{+1.69}_{-1.60}$ km s$^{-1}$), but they are within the respective

Table 5. Characteristics of our sample of hot BHB stars obtained from spectral simulations.

| Cluster/Star   | $\varv \sin i$ km s$^{-1}$ | $V_t$ km s$^{-1}$ | $\xi_{\text{max}}$ km s$^{-1}$ | $\xi = 0$ km s$^{-1}$ | [Fe/H] $\xi = 2$ km s$^{-1}$ | $n$ | $\xi = 0$ km s$^{-1}$ | [P/H] $\xi = 2$ km s$^{-1}$ | $n$ |
|---------------|--------------------------|------------------|-----------------------------|------------------------|--------------------------|-----|------------------------|--------------------------|-----|
| M 13/WF2-2541 | 2.4 $\pm$ 0.5            | -257.3 $\pm$ 0.5 | 0.0 $\pm$ 0.3               | -0.12 $\pm$ 0.20       | -0.19 $\pm$ 0.26         | 53  | +1.51 $\pm$ 0.11       | +1.46 $\pm$ 0.10         | 16  |
| M 13/WF4-3085 | 3.4 $\pm$ 0.8            | -255.4 $\pm$ 0.5 | 0.5 $\pm$ 0.3               | +0.72 $\pm$ 0.29       | +0.59 $\pm$ 0.37         | 117 | +1.84 $\pm$ 0.13       | +1.84 $\pm$ 0.15         | 12  |
| M 13/WF4-3485 | 3.0 $\pm$ 0.9            | -246.9 $\pm$ 0.9 | 1.5 $\pm$ 0.3               | -0.04 $\pm$ 0.50       | -0.17 $\pm$ 0.27         | 50  | ...                    | ...                      | ... |
| M 15/B84      | 4.0 $\pm$ 1.3            | -108.3 $\pm$ 0.4 | 0.0 $\pm$ 0.3               | -0.14 $\pm$ 0.10       | -0.35 $\pm$ 0.24         | 25  | ...                    | [Ti/H] ...                | ... |
| M 15/B267     | 6.2 $\pm$ 1.2            | -114.7 $\pm$ 0.7 | 0.5 $\pm$ 0.3               | +0.45 $\pm$ 0.37       | +0.21 $\pm$ 0.51         | 46  | +0.09 $\pm$ 0.20       | -0.08 $\pm$ 0.24         | 16  |
| M 15/B279     | 4.7 $\pm$ 1.6            | -104.5 $\pm$ 0.6 | 1.0 $\pm$ 0.3               | +0.42 $\pm$ 0.34       | +0.20 $\pm$ 0.40         | 71  | -0.44 $\pm$ 0.10       | -0.54 $\pm$ 0.11         | 11  |
In conclusion, the results shown here add to the mounting evidence of the existence of vertical abundance stratification, and hence atomic diffusion, in the atmospheres of BHB stars.

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Table 3. List of iron spectral lines used for the abundance analysis.

| \( \lambda, \text{Å} \) | \( \log gf \) | \( E_v, \text{cm}^{-1} \) | \( \log \gamma_{\text{rad}} \) | Object name |
|-----------------|---------------|-----------------|-----------------|-----------------|
| \( \text{Fe II} \) | WF2-2541 | WF4-3085 | WF2-3585 | B84 | B267 | B279 |
| 3903.756 | -1.555 | 60 807.23 | 8.99 | R&U | x | |
| 3906.035 | -1.704 | 44 929.55 | 8.72 | R&U | x | x | x |
| 3918.528 | -2.485 | 47 674.72 | 8.55 | R&U | x | x | x |
| 3920.635 | -1.335 | 60 625.45 | 9.01 | R&U | x | x | x |
| 3924.838 | -1.103 | 78 137.36 | 8.93 | R&U | x | |
| 3938.970 | -1.932 | 47 674.72 | 8.61 | R&U | x | x | |
| 3945.210 | -4.440 | 13 673.18 | 8.46 | R&U | x | |
| 4002.083 | -4.481 | 22 409.85 | 8.53 | R&U | x | |
| 4002.543 | -2.072 | 48 039.09 | 8.68 | R&U | x | |
| 4016.367 | -1.420 | 75 600.93 | 9.01 | R&U | x | |
| 4024.547 | -2.439 | 36 252.92 | 8.49 | R&U | x | x | x |
| 4032.945 | -1.920 | 79 331.50 | 9.18 | R&U | x | x | |
| 4044.012 | -2.671 | 44 929.55 | 8.70 | R&U | x | |
| 4048.832 | -2.381 | 44 917.07 | 8.73 | R&U | x | x | x |
| 4122.668 | -3.300 | 20 830.58 | 8.49 | R&U | x | x | x |
| 4128.748 | -3.578 | 20 830.58 | 8.61 | R&U | x | x | x |
| 4173.461 | -2.617 | 20 830.58 | 8.61 | R&U | x | x | x | x |
| 4177.618 | -3.776 | 73 395.93 | 8.61 | R&U | x | x | x |
| 4177.692 | -3.449 | 20 516.96 | 8.61 | R&U | x | x | x |
| 4180.862 | -2.535 | 20 830.58 | 8.49 | R&U | x | x | x | x |
| 4184.261 | -1.938 | 90 397.87 | 8.35 | R&U | x | |
| 4217.899 | -0.976 | 90 777.81 | 8.51 | VALD | x | |
| 4233.172 | -1.947 | 20 830.58 | 8.61 | R&U | x | x | x |
| 4258.154 | -3.478 | 21 812.06 | 8.48 | R&U | x | x | x |
| 4258.340 | -4.301 | 21 308.04 | 8.61 | R&U | x | x | x |
| 4263.869 | -1.696 | 62 049.03 | 8.95 | R&U | x | |
| 4273.326 | -3.303 | 21 812.06 | 8.62 | R&U | x | x | x | x |
| 4286.280 | -1.730 | 62 171.64 | 8.96 | R&U | x | x | x | x |
| 4296.572 | -2.933 | 21 812.06 | 8.49 | R&U | x | x | x | x |
| 4303.176 | -2.511 | 21 812.06 | 8.61 | R&U | x | x | x | x |
| 4314.979 | -3.193 | 38 164.19 | 8.49 | R&U | x | x | x | x |
| 4351.769 | -2.216 | 21 812.06 | 8.61 | R&U | x | x | x | x |
| 4357.584 | -2.010 | 49 100.98 | 8.64 | R&U | x | x | x | x |
| 4361.205 | -2.827 | 62 125.60 | 8.94 | R&U | x | x | x | x |
| 4361.247 | -2.264 | 49 506.93 | 8.70 | R&U | x | x | x | x |
| 4369.411 | -3.584 | 22 409.85 | 8.48 | R&U | x | x | x | x |
| 4384.319 | -3.684 | 21 430.36 | 8.48 | R&U | x | x | x | x |
| 4385.381 | -2.581 | 22 409.85 | 8.62 | R&U | x | x | x | x |
| 4416.830 | -2.602 | 22 409.85 | 8.61 | R&U | x | x | x | x | x |
| 4451.551 | -1.907 | 49 506.93 | 8.74 | R&U | x | x | x | x | x |
| 4455.266 | -2.000 | 50 212.83 | 8.58 | R&U | x | |
| 4461.694 | -4.002 | 63 272.98 | 9.02 | R&U | x | x | |
| 4461.706 | -2.065 | 50 212.83 | 8.62 | R&U | x | x | |
| 4472.929 | -3.531 | 22 939.36 | 8.48 | R&U | x | x | |
| 4489.183 | -2.971 | 22 810.36 | 8.49 | R&U | x | x | |
| 4491.405 | -2.755 | 23 031.30 | 8.48 | R&U | x | x | x | x | x |
| 4493.529 | -1.555 | 63 876.32 | 9.00 | R&U | x | |
Table 3, continued.

| \( \lambda, \) Å | \( \log g \) | \( E_\nu, \text{cm}^{-1} \) | \( \log \gamma_{\text{rad}} \) | Ref. | Object name |
|----------------|---------|----------------|----------------|-----|------------|
|                |         |                 |                |     | WF2-2541  |
| 4508.214       | −2.589  | 63 948.79       | 9.01           | R&U | x          |
| 4508.288       | −2.349  | 23 031.30       | 8.62           | R&U | x          |
| 4515.339       | −2.540  | 22 939.36       | 8.49           | R&U | x          |
| 4520.224       | −2.617  | 22 637.21       | 8.49           | R&U | x          |
| 4522.634       | −2.169  | 22 939.36       | 8.61           | R&U | x          |
| 4534.168       | −3.364  | 23 031.30       | 8.49           | R&U | x          |
| 4541.524       | −2.973  | 23 031.30       | 8.61           | R&U | x          |
| 4549.192       | −1.767  | 47 674.72       | 8.70           | R&U | x          |
| 4549.474       | −2.016  | 22 810.36       | 8.49           | R&U | x          |
| 4555.893       | −2.421  | 22 810.36       | 8.49           | R&U | x          |
| 4576.340       | −2.976  | 22 939.36       | 8.61           | R&U | x          |
| 4579.527       | −2.343  | 50 212.83       | 8.64           | R&U | x          |
| 4582.835       | −3.224  | 22 939.36       | 8.49           | R&U | x          |
| 4583.837       | −1.867  | 22 637.21       | 8.61           | R&U | x          |
| 4596.015       | −1.956  | 50 212.83       | 8.74           | R&U | x          |
| 4598.494       | −1.512  | 62 945.04       | 9.00           | R&U | x          |
| 4620.521       | −3.315  | 22 810.36       | 8.61           | R&U | x          |
| 4628.786       | −1.724  | 63 272.98       | 9.01           | R&U | x          |
| 4629.339       | −2.478  | 22 637.21       | 8.48           | R&U | x          |
| 4635.316       | −1.578  | 48 039.09       | 8.73           | R&U | x          |
| 4656.981       | −3.643  | 23 317.63       | 8.61           | R&U | x          |
| 4666.758       | −3.368  | 22 810.36       | 8.48           | R&U | x          |
| 4731.453       | −3.127  | 23 317.63       | 8.61           | R&U | x          |
| 4798.673       | −1.862  | 83 136.49       | 8.95           | R&U | x          |
| 4826.683       | −0.512  | 82 978.68       | 8.95           | R&U | x          |
| 4913.181       | −1.238  | 82 978.68       | 8.96           | R&U | x          |
| 4923.927       | −1.504  | 23 317.63       | 8.49           | R&U | x          |
| 4948.096       | −0.239  | 83 136.49       | 8.99           | R&U | x          |
| 4948.793       | −0.043  | 83 459.67       | 9.01           | R&U | x          |
| 4951.584       | 0.198   | 83 136.49       | 8.96           | R&U | x          |
| 4958.822       | −0.693  | 83 713.54       | 8.95           | R&U | x          |
| 4977.035       | −0.041  | 83 538.54       | 9.00           | R&U | x          |
| 4977.923       | −0.606  | 83 308.19       | 9.01           | R&U | x          |
| 4994.888       | 0.074   | 83 308.19       | 9.01           | R&U | x          |
| 5007.050       | 0.193   | 83 308.19       | 8.99           | R&U | x          |
| 5001.864       | 0.010   | 31 302.42       | 7.97           | R&U | x          |
| 5001.959       | 0.916   | 82 978.68       | 9.01           | R&U | x          |
| 5004.195       | 0.504   | 82 853.66       | 8.93           | R&U | x          |
| 5006.841       | −0.260  | 83 713.54       | 8.93           | R&U | x          |
| 5007.447       | −0.435  | 83 726.36       | 8.94           | R&U | x          |
| 5007.739       | −0.273  | 82 978.68       | 9.03           | R&U | x          |
| 5018.440       | −1.345  | 23 317.63       | 8.49           | R&U | x          |
| 5021.594       | −0.189  | 82 978.68       | 9.05           | R&U | x          |
| 5022.420       | −0.066  | 83 459.67       | 9.01           | R&U | x          |
| 5022.792       | −0.077  | 82 978.68       | 9.10           | R&U | x          |
| 5026.806       | −0.427  | 83 136.49       | 9.03           | R&U | x          |
Table 3. continued.

| $\lambda$, Å | $\log gf$ | $E_\text{i}$, cm$^{-1}$ | $\log \gamma_{\text{rad}}$ | Ref. | Object name |
|-------------|---------|----------------|-----------------|-----|-------------|
|            |         |             |                 |     | WF2-2541 | WF4-3085 | WF2-3585 | B84 | B267 | B279 |
| 5029.097   | −0.615  | 83 558.54   | 9.07            |     | x         |           |           |     |     |     |
| 5030.630   | 0.433   | 82 978.68   | 9.02            | R&U | x         | x         | x         |     |     |     |
| 5032.712   | 0.107   | 83 812.32   | 8.99            | R&U | x         | x         | x         |     |     |     |
| 5034.003   | −0.793  | 83 558.54   | 9.02            | R&U | x         |           |           |     |     |     |
| 5035.708   | 0.639   | 82 978.68   | 8.94            | R&U | x         | x         | x         | x   |     |     |
| 5045.114   | 0.005   | 83 136.49   | 9.05            | R&U | x         |           |           |     |     |     |
| 5047.641   | −0.216  | 83 136.49   | 9.04            | R&U | x         |           |           |     |     |     |
| 5061.718   | 0.290   | 83 136.49   | 9.06            | R&U | x         |           |           |     |     |     |
| 5067.893   | −0.063  | 83 308.19   | 9.06            | R&U | x         |           |           |     |     |     |
| 5070.899   | 0.275   | 83 136.49   | 9.02            | R&U | x         | x         | x         |     |     |     |
| 5075.764   | 0.168   | 84 326.91   | 8.94            | R&U | x         | x         | x         |     |     |     |
| 5082.230   | −0.121  | 83 990.06   | 9.00            | R&U | x         |           |           |     |     |     |
| 5089.214   | 0.025   | 83 308.19   | 9.05            | R&U | x         |           |           |     |     |     |
| 5097.271   | 0.315   | 83 713.54   | 8.94            | R&U | x         | x         | x         |     |     |     |
| 5100.607   | 0.234   | 83 726.36   | 8.95            | R&U | x         | x         | x         | x   |     |     |
| 5100.727   | 0.724   | 83 726.36   | 8.94            | R&U | x         | x         | x         | x   |     |     |
| 5106.097   | −0.961  | 83 812.32   | 9.00            | R&U | x         |           |           |     |     |     |
| 5106.109   | −0.234  | 83 308.19   | 9.06            | R&U | x         |           |           |     |     |     |
| 5117.034   | −0.033  | 84 131.56   | 9.01            | R&U | x         |           |           |     |     |     |
| 5123.188   | −2.457  | 62 125.60   | 8.49            | R&U | x         |           |           |     |     |     |
| 5123.188   | −2.286  | 84 131.56   | 8.99            | R&U | x         |           |           |     |     |     |
| 5132.669   | −0.094  | 82 637.21   | 8.54            | R&U | x         |           |           |     |     |     |
| 5143.880   | −0.223  | 84 266.56   | 9.00            | R&U | x         |           |           |     |     |     |
| 5149.465   | 0.544   | 84 266.56   | 8.92            | R&U | x         | x         |           |     |     |     |
| 5166.582   | −0.059  | 84 326.91   | 8.98            | R&U | x         |           |           |     |     |     |
| 5169.033   | −1.250  | 23 317.63   | 8.48            | R&U | x         | x         | x         | x   |     |     |
| 5180.314   | −0.103  | 83 812.32   | 9.00            | R&U | x         | x         |           |     |     |     |
| 5194.892   | −0.103  | 84 424.37   | 8.99            | R&U | x         |           |           |     |     |     |
| 5197.568   | −2.105  | 83 865.63   | 9.05            | VALD | x         | x         | x         | x   |     |     |
| 5197.577   | −2.348  | 26 055.42   | 8.48            | R&U | x         | x         | x         | x   |     |     |
| 5199.122   | 0.107   | 83 713.54   | 9.02            | R&U | x         | x         |           |     |     |     |
| 5213.960   | −0.264  | 84 527.78   | 9.01            | R&U | x         |           |           |     |     |     |
| 5214.057   | −0.633  | 83 713.54   | 9.04            | R&U | x         |           |           |     |     |     |
| 5215.844   | −0.153  | 83 726.36   | 9.22            | R&U | x         |           |           |     |     |     |
| 5216.854   | 0.480   | 84 710.68   | 8.96            | R&U | x         | x         | x         |     |     |     |
| 5216.863   | 0.659   | 84 527.78   | 8.93            | R&U | x         | x         | x         |     |     |     |
| 5218.842   | −0.181  | 83 726.36   | 9.10            | R&U | x         |           |           |     |     |     |
| 5223.260   | −0.190  | 83 812.32   | 9.04            | R&U | x         |           |           |     |     |     |
| 5223.290   | −0.914  | 84 266.56   | 9.09            | R&U | x         |           |           |     |     |     |
| 5223.800   | −0.532  | 83 713.54   | 9.01            | R&U | x         |           |           |     |     |     |
| 5224.411   | −0.450  | 83 990.06   | 9.09            | R&U | x         |           |           |     |     |     |
| 5227.323   | 0.191   | 84 444.83   | 8.93            | R&U | x         |           |           |     |     |     |
| 5227.481   | 0.834   | 84 296.83   | 8.91            | R&U | x         |           |           |     |     |     |
| 5232.787   | −0.115  | 83 726.36   | 8.93            | R&U | x         |           |           |     |     |     |
| 5234.615   | −2.385  | 84 266.56   | 8.98            | R&U | x         | x         | x         | x   |     |     |
| 5234.625   | −2.279  | 25 981.63   | 8.49            | R&U | x         | x         | x         | x   |     |     |
| 5237.951   | 0.101   | 84 266.56   | 9.11            | R&U | x         |           |           |     |     |     |
| λ, Å  | log $g_f$ | $E_i$, cm$^{-1}$ | log $\gamma_{rad}$ | Ref. | Object name |
|-------|----------|-----------------|---------------------|------|-------------|
|       |          |                 |                     |      |             |
| 5247.952 | 0.550 | 84 938.18 | 8.95 | R&U | x         |
| 5251.211 | −0.638 | 84 326.91 | 8.99 | R&U | x         |
| 5251.233 | 0.442 | 84 844.83 | 9.00 | R&U | x         |
| 5254.929 | −3.336 | 26 055.42 | 8.49 | R&U | x         |
| 5260.259 | 1.069 | 84 035.01 | 8.91 | R&U | x         |
| 5262.317 | −0.370 | 85 048.60 | 9.08 | R&U | x         |
| 5264.812 | −3.131 | 26 055.42 | 8.61 | R&U | x         |
| 5273.397 | −2.009 | 48 039.09 | 8.70 | R&U | x         |
| 5276.002 | −2.213 | 25 805.33 | 8.49 | R&U | x         |
| 5291.667 | 0.556 | 84 527.78 | 9.00 | R&U | x         |
| 5306.180 | 0.104 | 84 870.86 | 9.06 | R&U | x         |
| 5316.225 | 0.340 | 84 035.14 | 8.92 | R&U | x         |
| 5316.615 | −2.014 | 25 428.78 | 8.48 | R&U | x         |
| 5316.784 | −2.783 | 25 981.63 | 8.61 | R&U | x         |
| 5325.553 | −3.324 | 25 981.63 | 8.49 | R&U | x         |
| 5339.585 | 0.540 | 84 296.83 | 8.91 | R&U | x         |
| 5362.869 | −2.616 | 25 805.33 | 8.61 | R&U | x         |
| 5427.826 | −1.580 | 54 232.19 | 8.54 | R&U | x         |
| 5429.988 | 0.438 | 85 462.86 | 8.93 | R&U | x         |
| 5465.931 | 0.352 | 85 679.70 | 8.91 | R&U | x         |
| 5492.079 | 0.089 | 84 679.70 | 8.98 | R&U | x         |
| 5506.195 | 0.849 | 84 863.35 | 8.89 | R&U | x         |
| 5507.072 | −0.082 | 84 870.86 | 8.98 | R&U | x         |
| 5510.779 | −0.092 | 85 184.73 | 8.92 | R&U | x         |
| 5544.763 | 0.129 | 84 863.35 | 9.21 | R&U | x         |
| 5645.392 | 0.176 | 85 184.73 | 9.21 | R&U | x         |
| Ti II   |         |                 |                     |      |             |
| 3900.539 | −0.200 | 9118.28 | 8.36 | Pic. | x         |
| 3913.461 | −0.420 | 8997.78 | 8.37 | Pic. | x         |
| 4053.821 | −1.130 | 15 265.69 | 8.38 | Pic. | x         |
| 4163.644 | −0.130 | 20 891.78 | 8.37 | Pic. | x         |
| 4290.215 | −0.850 | 9395.79 | 8.46 | Pic. | x         |
| 4294.094 | −0.930 | 8744.34 | 8.22 | Pic. | x         |
| 4300.042 | −0.440 | 9518.14 | 8.47 | Pic. | x         |
| 4301.922 | −1.150 | 9363.73 | 8.47 | Pic. | x         |
| 4395.031 | −0.540 | 8744.34 | 8.21 | Pic. | x         |
| 4399.765 | −1.190 | 9975.99 | 8.46 | Pic. | x         |
| 4443.801 | −0.720 | 8710.56 | 8.20 | Pic. | x         |
| 4468.507 | −0.600 | 9122.14 | 8.21 | VALD | x         |
| 4501.270 | −0.684 | 8997.78 | 8.20 | VALD | x         |
| 4563.757 | −0.690 | 9851.00 | 8.22 | Pic. | x         |
| 4571.971 | −0.230 | 12 677.10 | 8.37 | Pic. | x         |
| 4805.085 | −0.960 | 16 623.11 | 8.26 | VALD | x         |
| 4911.195 | −0.610 | 25 192.96 | 8.30 | Pic. | x         |
| 5188.687 | −1.050 | 12 758.26 | 8.22 | Pic. | x         |
Table 3. continued.

| λ, Å | log gf | \( E_i \), cm\(^{-1} \) | \log \gamma_{rad} | Ref. | Object name |
|------|-------|---------------------|------------------|-------|-------------|
|      |       |                     |                  |       |             |
| 4178.463 | -0.409 | 77 671.28 | VALD | x | x |
| 4420.712 | -0.478 | 88 893.22 | VALD | x | x |
| 4463.027 | 0.026 | 105 553.90 | VALD | x | x |
| 4468.000 | -0.208 | 105 223.22 | VALD | x | x |
| 4475.270 | 0.301 | 105 553.90 | VALD | x | x |
| 4499.230 | 0.377 | 107 922.93 | VALD | x | x |
| 4530.823 | -0.020 | 105 303.87 | VALD | x | x |
| 4588.032 | 0.575 | 103 343.94 | VALD | x | x |
| 4589.846 | 0.500 | 103 166.50 | VALD | x | x |
| 4602.069 | 0.799 | 103 667.86 | VALD | x | x |
| 4658.309 | -0.412 | 103 665.56 | VALD | x | x |
| 4943.497 | 0.083 | 103 665.56 | VALD | x | x |
| 5253.479 | 0.291 | 88 893.22 | VALD | x | x |
| 5296.077 | -0.134 | 87 124.60 | VALD | x | x |
| 5316.055 | -0.341 | 86 743.96 | VALD | x | x |
| 5344.729 | -0.329 | 86 597.55 | VALD | x | x |
| 5425.880 | 0.241 | 87 124.11 | VALD | x | x |
| 5499.697 | -0.495 | 87 124.11 | VALD | x | x |
| 6024.178 | 0.137 | 86 745.03 | VALD | x | x |
| 6034.039 | -0.209 | 86 599.85 | VALD | x | x |
| 6043.084 | 0.384 | 87 124.11 | VALD | x | x |
| 6165.598 | -0.469 | 87 124.11 | VALD | x | x |

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VALD 1999POBeo..65..223K Kupka & Ryabchikova
Pic. 2001ApJS...132..403 Pickering, Thorne & Perez
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