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Research Article

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Assessing of channel structure and magnetic properties on heavy metal ions removal from water

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Abstract:

The synthesis of Li_{1.1}Co_{0.3}Fe_{2.1}O_{4} ferrite nanoparticles has been synthesized by the citrate auto combustion method. The distribution of cations on A-site and B-site was studied by X-ray Diffraction (XRD), High-Resolution Transmission Electron Microscopy (HRTEM). The crystallite size and strain were calculated by using the W–H (Williamson-Hall) plot. The crystallite of Li_{1.1}Co_{0.3}Fe_{2.1}O_{4} ferrite nanoparticle is ∼30 nm. HRTEM confirmed the homogeneous formation of the cubic phase. The calculated height and spacing parameters related to roughness are essential to achieve the efficiency of Li_{1.1}Co_{0.3}Fe_{2.1}O_{4} to be used in micro-batteries, smart windows, smart mirrors, displays, gas sensors, and other applications. According to the obtained data, the Li_{1.1}Co_{0.3}Fe_{2.1}O_{4} has a spiky surface with R_{ku} = 5.50. Additionally, the magnetic hysteresis loop has been clarified using the Vibrating Sample Magnetometer (VSM). The double peak characteristic in the Switching field distribution (SFD) reveals the competition between exchange coupling and strong dipolar interactions. Li_{1.1}Co_{0.3}Fe_{2.1}O_{4} has employed as a sorbent material for the removal of lead (II) ions from wastewater. The main advantages of the synthesized sample are ease of separation, high adsorption, low cost as well as recycled with notable efficiency. Two models of adsorption isotherms (Freundlich and Langmuir) are utilized to recognize the adsorption mechanism.

Keywords: isothermal models, adsorption process, lead removal, Li nano ferrite, spacing parameters.

1. Introduction

A great properties of spinel ferrites are inexpensive, abundant in nature, environment-friendly and more stable.

Lithium ferrite in the spinel structure is a low-cost material that, is generally valuable for a number of unique electromagnetic properties of that enlarges the scope their application in various devices of modern microwave technology[1-2]. They also can be used as cathode materials in lithium-ion batteries [3], and as sensing units in gas sensors [4]. It has a high-resistivity, low mobility semiconductor, and low eddy current losses. The properties of lithium ferrite can be enhanced by the substitution of divalent (cobalt) ions.

The spinel cobalt ferrite (CoFe_{2}O_{4}) materials have great attention due to their special physical properties.

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and catalytic properties. It was studied in several fields of electrocatalysts [5], chemical sensors[6-7], Li-ion batteries [8-9], and supercapacitors [10-11]. The structure of CoFe₂O₄ has an inverse spinel structure, signified as (Fe³⁺)[Co²⁺Fe³⁺]O₄, where Fe³⁺ ions occupied one-eighth of the tetrahedral sites (A sites) and Co²⁺ and Fe³⁺ ions occupied one-half of the octahedral sites (B sites) ions [12-13]. This structure has revealed good magnetic, optical, and electrical conductivity [14-15].

Removal of Heavy Metals from wastewater is a significant issue because they not only contaminate the water bodies but are also toxic to the ecosystem. As the majority of the heavy metals are non-degradability and highly toxic in nature. Consequently, their concentrations have to be decreased to appropriate levels before discharging into the environment, or else these can cause a threat to human as well as animal health.

In the present study, Li₁.₁Co₀.₃Fe₂.₁O₄ is synthesized using the citrate precursor method because of its cost - effectiveness, homogeneous formation of cubic ferrites, less time-consuming technique, and uniform particle size. Structural, thermal analysis morphological, and magnetic investigations were performed on the investigated sample. Further, the sample was employed as a sorbent for the removal of heavy metal (lead II) from water. The effect of pH on the adsorption process has been studied.

2. Experimental Work

2.1 Preparation technique

All chemicals were purchased from Sigma-Aldrich. LiCoFeO was prepared by mixing non-stoichiometric proportions of iron, cobalt, lithium nitrates with a calculated amount of citric acid by citrate auto combustion technique as reported previously [16] with some adjustments as shown in Fig. 1a

2.2 Adsorption test

The batch equilibrium technique was employed as mentioned in the previous work [17]. The Pb²⁺ ion concentration was measured by atomic absorption spectrophotometry ICP (spectrometry Prodigy7).
Fig. 1: a-b Flowchart for the synthetic scheme and the recovery of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ from water by a small magnet

The adsorption efficiency ($\eta$) and the equilibrium adsorption capacity ($q$) of metal ions were calculated according to the following equations [18]

$$\eta = \frac{C_i - C_e}{C_i} \times 100$$  \hspace{1cm} (1)

$$q = \frac{(C_i - C_e)V}{m}$$  \hspace{1cm} (2)

Where $C_i$ and $C_e$ are the initial and equilibrium (final) concentrations (mg/L) of metal ion solution, respectively. While V and m are the volume of Pb (II) solution, and the mass of adsorbent respectively. The nano ferrite sample is attracted to the magnet and accordingly can be easily separated as illustrated in Fig. 1b.

3. Results and discussion:

Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ exhibits a non-stoichiometry. The non-stoichiometry may occur for some ferrite materials in which two valence states exist for one of the ion types. It is due to an excess positive charge which must be offset by some types of defects. This may be accomplished by the oxidation and reduction processes of cations like Fe$^{3+}$ and Fe$^{2+}$ during the auto combustion. In fact, the chemical formula is often written as Me$_{1-x}$M$_x$Fe$_2$O$_4$ to indicate a condition of non-stoichiometry with an excess of cations.
Figure 2a illustrates the XRD pattern of the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ sample revealing the formation of cubic spinel phase with space group (Fd-3m). The XRD data is compared with JCPDS card 04-022-8066. There are no appreciable impurity peaks in the spectrum of the sample. The indexed peaks are (220), (311), (222), (400), (422), (511), (440), (620) and (533). The more intense peak is observed at $\theta = 35.73^{\circ}$ with hkl (311) which agrees well with the reported literature [19].

Certainly, the stoichiometric ratio for the tetra/octa is 0.50 causes a single spinel phase. In the present study, this ratio equals 0.66 yielding an excess amount of metals. In fact, the XRD pattern does not detect any secondary phases. This can be attributed to the large surface-to-volume ratio leading to the flexibility of the structure as well as the flexibility of the composition [20]. Li-ferrite nanoparticles show a partially inverted spinel structure [21]. The incorporation of the two cations at both cationic lattice sites of the spinel structure allows a broad compositional range of stability, even with rich compositions (cations in A/cations in B >0.5). Additionally, this non-stoichiometry can be structurally accommodated by the formation of point defects.

The crystallite size ($D_{XRD}$) and the average lattice strain ($\varepsilon$) of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ can be detected according to the Williamson–Hall equations [22]:

$$\beta \cos \theta = 4\varepsilon \sin \theta + \frac{K\lambda}{D}$$

(b)

The XRD diffraction pattern and The Williamson–Hall plot for Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$
where $\beta$ is the FWHM of XRD peak; $\theta$ is the Bragg angle; K is the dimensionless shape factor ($K = 0.9$); $\lambda$ is the wavelength of CuK$\alpha$ radiation ($\lambda = 1.5406$). Figure 2b shows the fitting plots of $\beta \cos \theta$ versus $4 \sin \theta$, and the $\epsilon$ is determined from the slope. The calculated values of $\epsilon$ and $D_{\text{XRD}}$ are listed in Table 1. The strain is considered to be uniform in all the crystallographic directions. The W-H plot for the investigated sample shows the negative value of the strain indicates the compressive nature of it which leads to lattice shrinkage.

The dislocation density $\delta$, represents the total defects in the sample. It is defined as the length of dislocation lines per unit volume of the crystal and it is calculated using the following equation [23]:

$$\delta = \frac{1}{D^2} \quad (4)$$

Other important parameters that reveal the crystallization of the sample crystals are the distortion parameter ($g$), inter-chain distance ($r$), and theoretical density ($D_x$). These parameters can be calculated from XRD data by using the following equations [24][25];

$$g = \frac{\beta}{\tan \theta} \quad (5)$$

$$r = \frac{5\lambda}{8 \sin \theta} \quad (6)$$

$$D_x = \frac{ZM}{NV} \quad (7)$$

where $Z$ (8) is the number of molecules per unit cell, $M$ is the molecular weight, $N$ is Avogadro's number and $V$ is the unit cell volume. All the calculated parameters are tabulated in the Table.

| Sample         | (a) °A | D (nm) | $\delta$ (nm$^{-2}$) | (g) | $\epsilon \times 10^{-5}$ | r (nm) | V (nm$^3$) | Dx (g cm$^{-3}$) |
|----------------|-------|-------|----------------------|-----|--------------------------|--------|-----------|------------------|
| Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ | 8.324 | 30.7  | 0.001                | 0.013 | 43                       | 0.313  | 0.576      | 4.758            |

The prepared sample is of high quality with good crystallinity and fine crystallite size as confirmed from the calculated parameters in the Table.

Fig. 3a-b illustrates the high-resolution transmission electron micrographs (HRTEM) for Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ sample, inset (a) the selected area electron diffraction patterns (SAED) and inset (b) histogram of the size distribution.
The shape of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nanoparticle is roughly spherical and highly agglomerated due to the magnetic nature of the sample. The selected area electron diffraction (SAED) consists of concentric rings with spots over the rings. This characteristic specifies that the sample is good nanocrystalline in nature [26]. The rings with a dotty pattern in SAED ratify the varied size distribution of the nano ferrite sample. The inset in Fig. 3b shows the size distribution of the sample and the mean particle size is 70 nm.

Atomic force microscopy (AFM) is a crucial tool to scrutinize the surface topography of the material. Fig. 4 shows the 2D - AFM image of the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nano ferrite particle. The images show that Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ exhibits waviness surface texture. The hill region is comparatively rough compared to the valley region. In other words, the typical roughness is the same for surfaces with roughness profiles completely different, because it takes into consideration the average profile of heights only. For this reason, additional sophisticated parameters can be clarified to fully describe a surface when more significant data is required.

The profile peak height ($R_p$) is the measure of the highest peak around the surface profile from the baseline. Similarly, the valley depth ($R_v$) is the measure of the deepest valley across the surface profile analyzed from the baseline. Consequently, the maximum height of the profile ($R_T$) can be identified as the vertical distance between the minimum valley and the maximum peak as shown in Fig. 4c.

$$R_T = R_p + R_v \quad (8)$$
Table 2 show Maximum Profile Valley Depth $R_v$, Maximum Profile Peak Height $R_p$, Maximum Height of the Profile $R_t$, average roughness $R_a$, root mean square roughness $R_q$, skewness of the line $R_{sk}$, and kurtosis of the line $R_{ku}$.

| Sample          | $R_v$ nm | $R_p$ nm | $R_t$ nm | $R_a$ nm | $R_q$ nm | $R_{sk}$ | $R_{ku}$ |
|-----------------|----------|----------|----------|----------|----------|----------|----------|
| Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ | 5.03     | 6.73     | 11.8     | 0.707    | 1.06     | 1.0      | 5.50     |

The roughness parameters are estimated by analyzing the topography scans of the sample’s surface and the obtained data is tabulated in Table 2. As shown from the Table, the surface profile parameters include average roughness ($R_a$), root mean square roughness ($R_q$), maximum peak to valley height ($R_t$), skewness of the line ($R_{sk}$), kurtosis of the line ($R_{ku}$). These parameters are important for obtaining specified data, like a scratch or an unusual crack on the material.

Fig. 4 a-c  AFM image of the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nano ferrite sample.
Roughness skewness ($R_{sk}$) is used to measure the symmetry and is more sensitive to occasional deep valleys or high peaks. Usually, $R_{sk}$ is used to distinguish two profiles of the same $R_a$ or $R_q$ values but of different shapes.

Roughness kurtosis ($R_{ku}$) is utilized to estimate the distribution of the spikes above and below the mean plane. For spiky surfaces, $R_{ku} > 3$; for bumpy surfaces, $R_{ku} < 3$. According to the value of $R_{ku}$ the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ has a spiky surface ($R_{ku} = 5.50$).

Finally, roughness is a very significant parameter for numerous applications. The characterization of materials through their roughness allows one to attain data on the efficiency of samples in many application areas. The height and spacing parameters associated with roughness are essential to achieve the efficiency of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ to be used in micro-batteries, smart mirrors, gas sensors, and other applications.

The adsorption of molecules on surfaces as a function of the relative pressures ($P/P_0$) of the adsorbate is elucidated by Brunauer, Emmett, and Teller theory.

![Fig. 5 a-b Type of sorption isotherms and pores distribution of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$.](image)

Figure 5a illustrates the adsorption/desorption isotherm of nitrogen at 77 K for Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ ferrite nanoparticles. As shown from the plot that the desorption branch does not follow the
adsorption branch. It forms an obvious hysteresis loop. This loop is due to the amount of adsorbed along the desorption branch and the adsorption branch is quite different. The sorption isotherm is categorized as type IV and the hysteresis loop is of type H1. This agrees well with the IUPAC (International Union of Pure and Applied Chemistry) classification.

The distribution of pore radius versus pore volume of the studied sample is shown in Fig. 5b. Adsorption is dependent on the surface area, pore sizes as well as pore volume. The BET surface area of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ is about 127.14 m$^2$/g. The pore volume and pore sizes are 0.0436 cm$^3$/g and 2.034 nm respectively. According to the value of the pore size, the investigated sample can be classified as mesoporous materials [27]. In mesoporous materials, the sorption efficiency depends on two issues the first one is fluid wall interaction strength. The second is the attractive interactions between the fluid molecules. Consequently, multilayer adsorption and capillary condensation can occur in the pore.

Measurement of magnetization at room temperature for Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nanocrystalline spinel powder is carried out using the VSM technique. Figure 6a-b illustrates the magnetization versus external magnetic field (M–H) hysteresis loop and the fitted Langevin curve. The experimental M versus H loop is determined with an external applied magnetic field ranging between $-20$ and $+20$ kOe. The saturation magnetization value that corresponds to 20 kOe is 39.898 emu/g. The saturation magnetization (Ms) of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nanoparticle also can be detected by an approximation of Stoner–Wohlfarth theory by extrapolating the plot of magnetization versus $1/H^2$ to approach zero [28–30]. This theory is applied for non-interacting particles as confirm from squareness values in the Tables. In this way, the Ms value is equal to 40.10 emu/g. The obtained value is very comparable to the experimental value signifying that an applied field of ±20 kOe is appropriate to saturate the investigated samples.

The Li$^{1+}$ ion is a nonmagnetic cation (0 µB) that does not respond to the magnetization of the sub-lattice. Consequently, it does not impart to the net magnetic moment. Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ ferrite has a partially inverse structure so most of the lithium ions occupy octahedral sites (B-sites) where as Fe$^{3+}$ (5 µB) and Co$^{2+}$ (3 µB) ions are distributed on both octahedral (B-sites) and tetrahedral sites (A-sites) [31,32]. Moreover, cobalt reveals an irregular trend when incorporated in Li ferrites.
The magnetization trend of the ferrite samples is due to the contributions of the magnetic moments from Fe$^{3+}$ and Co$^{2+}$ at B-sites [33,34]. Neel’s phenomenological model is applied to elucidate the magnetic parameters in LiCoFeO$_4$ ferrites nanoparticles. The obtained value of saturation magnetization is lower than that found in previous studies of Li$_{0.5}$Fe$_{2.5}$O$_4$ nanoparticles [35]. The less magnetization of the samples is due to the fact that the magnetization of A-sublattice becomes much diluted and the A-B exchange interactions become weaker or comparable with the B-B exchange interactions. Consequently, the canted spins and the Yafet-Kittel (Y-K) angle are increased [36]. Additionally, it can be also owed to the formation of an inactive magnetic layer [37,38] and the disordered cation distribution [39].

From the hysteresis loops, the coercivity ($H_c$), saturation magnetization ($M_s$), remnant magnetization ($M_r$), squareness ($M_r/M_s$), exchange bias [40], and magnetic moment ($n_B$) [41] are calculated and tabulated in Table 3.

Coercivity is often considered an extrinsic property of materials. It is sensitive to defects as grain boundaries, and dislocations. The value of the squareness ratio ($M_r/M_s$) is an indication of the soft character of the nanoparticles and its value varies from 0 to 1 [42]. In the present case, the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ sample interacts magneto statically according to Stoner–Wolfarth model [46].
The anisotropy constant is determined from the Stoner-Wohlfarth equation as follows

\[ K = \frac{H_c \times M_S}{0.96} \]  

(9)

Where \( H_c \) is the coercivity, \( M_s \) saturation magnetization and \( K \) is magnetic anisotropy constant. The weak anisotropy of the investigated samples is due to the non-collinearity (canting) of spins on their surface.

A noticeable shift of the loop is detected and is ascribed to exchange bias phenomena. The exchange bias field are detected from the following equation and tabulated in the Table.

\[ H_{\text{exch}} = -\frac{(H_{\text{left}} + H_{\text{right}})}{2} \]  

(10)

This is referred to the existence of different spin configurations. In nanostructure sample, the magnetic structure of the core is different from that of the surface where spin frustration predominates [47,48].

| Sample          | \( M_s \) (emu.g\(^{-1}\)) | \( M_r \) (emu.g\(^{-1}\)) | \( H_c \) (Oe) | M-H loop area (erg.g\(^{-1}\) x10\(^3\)) | \( K \) (emu.Oe.g\(^{-1}\)) | Squareness | \( n_B \) (\( \mu_B \)) | Exchange bias |  |
|-----------------|-----------------------------|-----------------------------|---------------|------------------------------------------|-----------------------------|-------------|-----------------|-----------------|---|
| Li\(_{1.1}\)Co\(_{0.3}\)Fe\(_{2.1}\)O\(_4\) | 39.898                      | 12.472                      | 247.81        | 44.195                                   | 9,887                       | 0.307       | 1.475           | 13.3            |   |

The Switching Field Distribution (SFD) for the materials is one of the significant criteria for magnetic recording. It has a high impact on high-density recording performance [49].

Fig. 7 shows the SFD value for the Li\(_{1.1}\)Co\(_{0.3}\)Fe\(_{2.1}\)O\(_4\) as a function of coercivity. Here, the SFD value is calculated from eq. (11) as the ratio of the half-width of the peak in the derivative M-H and coercivity

\[ \text{SFD} = \frac{\Delta H}{H_c} \]  

(11)

The derivative curve shows a double peak typical of a two-step reversal. Generally, the double peak characteristic in the derivative reveals the competition between exchange coupling and strong dipolar interactions.
Tuning the dipolar field energy to a value of the same order as the exchange coupling energy leads to a double peak derivative whose first peak usually occurs before reaching zero field coming from the saturation field [50-52]. Therefore, in our sample, the dipolar interactions are more effective, which implies a strong reduction of the inter bump exchange coupling in our bump array.

![Graph](image_url)

**Fig. 7** SFD value for the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ as a function of coercivity.

From the above arguments, it is concluded that the small SFD value (0.544) of the investigated sample is due to an intrinsic magnetic nature such as uniformity of substitution and crystallinity as confirmed from XRD data.

The removal of Pb II heavy metal cations from an aqueous solution is detected through batch experiments using nano-Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ as an adsorbent. The interaction between adsorbent and adsorbate can be adjusted by experimental parameters as pH values [51]. The concentration of heavy metal nitrates is measured by the ICP technique. The removal efficiency of investigated sample is determined using Equations (1,2) as mentioned in the experimental section.

Figure 8 illustrates the effect of pH on the adsorption efficiency of Pb$^{2+}$ ions in the aqueous solution. The pH can influence the sorption capacity of Pb$^{2+}$ ions by changing the surface charges of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$, the Pb$^{2+}$ complex type, and the number of active sites available for the adsorption process [52].
The results indicate that Pb$^{2+}$ removal increases to the maximum with increasing pH from 4 to 8 at room temperature (30 °C) as shown in Table (4). In the present study, the maximum efficiency of Pb$^{2+}$ removal is about 99.82 % at pH 8. This is elucidated in view of a competition between H$^+$ and Pb (II) ions to be adsorbed on Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$. At lower pH values, H$^+$ ions are able to prevent a significant number of adsorption sites at Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ from the Pb$^{2+}$ adsorption process [52]. While, at high pH values, the H$^+$ ions competition nearly vanishes and the positively charged Pb$^{2+}$ and Pb (OH)$^+$ ions can be attached to the free binding sites, increasing the uptake of the Pb$^{2+}$ ions [53].

Table 4 and Fig. (8) Effect of pH on the adsorption efficiency of Pb$^{2+}$ ions.

| Pollutant | PH | $C_i$ (mg.L$^{-1}$) | $C_e$ (mg.L$^{-1}$) | $\eta$ % |
|-----------|----|---------------------|---------------------|---------|
| Pb        | 4  | 6.501               | 3.018               | 53.56   |
|           | 5  | 2.998               | 0.065               | 99.00   |
|           | 6  | 0.022               | 0.012               | 99.82   |

The most recurrently utilized models to designate the adsorption results are the Freundlich and Langmuir isotherms. In the case study, these models are applied to scrutinize the adsorption of the heavy metals on Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ at different Pb$^{2+}$ ion concentrations. This study is carried out for equilibrium metal ion concentrations varying from 11.9351 to 0 mg/L. The constants of the two models are calculated using the following relations and the results are presented in the inset of Fig. 9a-b. The Langmuir equation is given by [17]:

$$\frac{C_e}{q_e} = \frac{1}{q_m K_L} + \frac{C_e}{q_m}$$  \hspace{1cm} (12)

While the Freundlich logarithmic form of equation is specified as [54]

$$Ln q_e = Ln K_f + \frac{1}{n} Ln C_e$$  \hspace{1cm} (13)

Where $C_e$, equilibrium metal concentration (mg/L), $q_m$ and $K_L$ are the Langmuir constants accompanying with maximum adsorption capacity (mg/g), and the relative energy of adsorption
(L/mg), respectively. $q_e$ is the equilibrium sorption capacity (mg/g), $K_F$ and $n$ are Freundlich equilibrium constants.

It indicates that the experimental data fitted well to both isotherm models. The Langmuir model, considering a monolayer adsorption process of molecules on solid surfaces, gives a good model for the adsorption system. According to this model, see Eq. (12) and Fig. 9a, the maximum sorption capacity, $q_m$, is estimated as 102.04 mg/g. The Langmuir constant, $K_L$, equals 0.2503 L/mg which designates high sorption energy between Pb II and Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$. The electrostatic nature of bonding between Pb$^{2+}$ and Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nano ferrite sample is the main reason for the observed high sorption of Pb$^{2+}$ ions on Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$.

Furthermore, the parameters ($K_f$ and 1/$n$) calculated from the intercept and slope of this Fig. 9b, respectively are given in the inset of the figure. The small value of 1/$n$ (less than 1) for the adsorption of lead Pb$^{2+}$ on the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ showing that the adsorption process is favorable.

**Conclusion**

Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nano sample was prepared in single phase cubic spinel with space group (Fd-3m). XRD assure that the sample was prepared in nano scale with crystallite size 30 nm. The shape of nanoparticles is roughly spherical and highly agglomerated. M/M$_s$ ratio assure the sample Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ has magneto static interaction. The sample has a weak ferrimagnetic
behavior with Curie temperature near 715K. The adsorption data are fitted to the Freundlich and Langmuir isotherm models. The best fitting is obtained by utilizing the Freundlich model for Pb $^{+2}$ metal i.e. PbII with $R^2$ value 0.8835.

Declarations

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Conflict of interest

The authors declare that they have no conflict of interest.

 Availability of data and material

Not applicable

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References

1. R. Cheruku, G. Govindaraj, and L. Vijayan, Mater. Chem. Phys. 146, 389 (2014).
2. Y.-P. Fu and S.-H. Hu, Ceram. Int. 36, 1311 (2010).
3. H. Zeng, T. Tao, Y. Wu, W. Qi, C. Kuang, S. Zhou, and Y. Chen, RSC Adv. 4, 23145 (2014).
4. N. Rezlescu, C. Doroftei, E. Rezlescu, and P. D. Popa, Sensors Actuators B Chem. 133, 420 (2008).
5. H. Zhu, S. Zhang, Y.-X. Huang, L. Wu, and S. Sun, Nano Lett. 13, 2947 (2013).
6. R. Yadav, M. K. Yadav, and N. K. Singh, Int. J. Electrochem. Sci 8, 6321 (2013).
7. X. Zhao, Y. Fu, J. Wang, Y. Xu, J.-H. Tian, and R. Yang, Electrochim. Acta 201, 172 (2016).
8. C. H. Kim, Y. Myung, Y. J. Cho, H. S. Kim, S.-H. Park, J. Park, J.-Y. Kim, and B. Kim, J. Phys. Chem. C 113, 7085 (2009).
9. Y. Wang, D. Su, A. Ung, J. Ahn, and G. Wang, Nanotechnology 23, 55402 (2012).
10. Z. Zhang, Y. Wang, M. Zhang, Q. Tan, X. Lv, Z. Zhong, and F. Su, J. Mater. Chem. A 1, 7444 (2013).
11. P. He, K. Yang, W. Wang, F. Dong, L. Du, and Y. Deng, Russ. J. Electrochem. 49, 359 (2013).
12. L. Lv, Q. Xu, R. Ding, L. Qi, and H. Wang, Mater. Lett. 111, 35 (2013).
13. G. A. Sawatzky, F. Van Der Woude, and A. H. Morrish, J. Appl. Phys. 39, 1204 (1968).
14. V. Neburchilov, H. Wang, J. J. Martin, and W. Qu, J. Power Sources 195, 1271 (2010).
15. E. Rios, J.-L. Gautier, G. Poillerat, and P. Chartier, Electrochim. Acta 44, 1491 (1998).
16. E. E. Ateia, M. K. Abdelmaksoud, and H. Ismail, J. Mater. Sci. Mater. Electron. 1 (2021).
17. E. E. Ateia, R. Ramadan, and A. S. Shafaay, Appl. Phys. A 126, 1 (2020).
18. R. Ramadan, S. I. El-Dek, and M. M. Arman, Appl. Phys. A 126, 1 (2020).
19. C. Singh, E. E. Ateia, S. Bindra Narang, M. Farag, J. Singh, D.E. El-Nashar, Ceramic International, 47,5 (2021)
20. K. Tanida and T. Kitamura, Tohoku Daigaku Senko Seiren Kenkyusho Iho 40, 71 (1984).
21. S. K. Gurav, S. E. Shirsath, R. H. Kadam, S. M. Patange, K. S. Lohar, and D. R. Mane, Mater. Res. Bull. 48, 3530 (2013).
22. V. D. Mote, Y. Purushotham, and B. N. Dole, J. Theor. Appl. Phys. 6, 1 (2012).
23. M. A. Amer, A. Matsuda, G. Kawamura, R. El-Shater, T. Meaz, and F. Fakhry, J. Alloys Compd. 735, 975 (2018).
24. E. E. Ateia, H. Ismail, H. Elshimy, and M. K. Abdelmaksoud, J. Inorg. Organomet. Polym. Mater. 1 (2021).
25. M. A. Amer, T. M. Meaz, S. S. Attalah, and A. I. Ghoneim, J. Alloys Compd. 654, 45 (2016).
26. S. L. Reddy, T. R. Reddy, N. Roy, R. Philip, O. A. Montero, T. Endo, and R. L. Frost, Spectrochim. Acta Part A Mol. Biomol. Spectrosc. 127, 361 (2014).
27. R. Schmidt, E. W. Hansen, M. Stoecker, D. Akpioriaye, and O. H. Ellestad, J. Am. Chem. Soc. 117, 4049 (1995).
28. A. Baykal, S. Esir, A. Demir, and S. Güner, Ceram. Int. 41, 231 (2015).
29. M. A. Almessiere, Y. Slimani, and A. Baykal, Ceram. Int. 44, 9000 (2018).
30. S. Güner, M. Amir, M. Geleri, M. Sertkol, and A. Baykal, Ceram. Int. 41, 10915 (2015).
31. A. Goldman, (n.d.).
32. I. Soibam, S. Phanjoubam, and C. Prakash, J. Magn. Magn. Mater. 321, 2779 (2009).
33. I. Soibam, S. Phanjoubam, and C. Prakash, J. Alloys Compd. 475, 328 (2009).
34. S. Ouaissa, A. Benyoussef, G. S. Abo, M. Ouaissa, and M. Hafid, Phys. Procedia 75, 792 (2015).
35. R. P. Patil, S. B. Patil, B. V Jadhav, S. D. Delekar, and P. P. Hankare, J. Magn. Magn. Mater. **401**, 870 (2016).

36. R. Topkaya, A. Baykal, and A. Demir, J. Nanoparticle Res. **15**, 1 (2013).

37. S. Rasheed, R. A. Khan, F. Shah, B. Ismail, J. Nisar, S. M. Shah, A. Rahim, and A. R. Khan, J. Magn. Magn. Mater. **471**, 236 (2019).

38. M. D. D. HOSSAIN, (2015).

39. T. Zeeshan, S. Anjum, S. Waseem, M. Riaz, and R. Zia, Ceram. Int. **46**, 3935 (2020).

40. C. Vázquez-Vázquez, M. A. López-Quintela, M. C. Buján-Núñez, and J. Rivas, J. Nanoparticle Res. **13**, 1663 (2011).

41. S. Panchal, S. Raghuvasanshi, K. Gehlot, F. Mazaleyrat, and S. N. Kane, AIP Adv. **6**, 55930 (2016).

42. U. Kurtan, R. Topkaya, A. Baykal, and M. S. Toprak, Ceram. Int. **39**, 6551 (2013).

43. B. K. Rai, L. Wang, S. R. Mishra, V. V Nguyen, and J. P. Liu, J. Nanosci. Nanotechnol. **14**, 5272 (2014).

44. G.-Y. Kim, J.-H. Jeon, M.-H. Kim, D. Suvorov, and S.-Y. Choi, Met. Mater. Int. **19**, 1209 (2013).

45. S. Chakrabarty, A. Dutta, and M. Pal, J. Alloys Compd. **625**, 216 (2015).

46. E. E. Ateia and F. S. Soliman, Mater. Sci. Eng. B **244**, 29 (2019).

47. S. R. Naik, A. V Salker, S. M. Yusuf, and S. S. Meena, J. Alloys Compd. **566**, 54 (2013).

48. K. Kamazawa, Y. Tsunoda, H. Kadowaki, and K. Kohn, Phys. Rev. B **68**, 24412 (2003).

49. O. Hellwig, A. Berger, T. Thomson, E. Dobisz, Z. Z. Bandic, H. Yang, D. S. Kercher, and E. E. Fullerton, Appl. Phys. Lett. **90**, 162516 (2007).

50. T. Hauet, L. Piraux, S. K. Srivastava, V. A. Antohe, D. Lacour, M. Hehn, F. Montaigne, J. Schwenk, M. A. Marioni, and H. J. Hug, Phys. Rev. B **89**, 174421 (2014).

51. O. E. A. Salam, N. A. Reid, and M. M. ElShafei, J. Adv. Res. **2**, 297 (2011).

52. M. Horsfall Jrn and A. I. Spiff, Electron. J. Biotechnol. **7**, 14 (2004).

53. H. Liu, *Mechanisms of the Oxidation of Lead (II) Solid Phases by Chlorine and Its Relationships with Colloidal Lead Release* (University of Washington, 2010).
Figures

**Figure 1**

a-b Flowchart for the synthetic scheme and the recovery of Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ from water by a small magnet

![Flowchart](image1)

**Figure 2**

a-b The XRD diffraction pattern and The Williamson–Hall plot for Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$

![XRD Pattern](image2)
Figure 3

(a-b) HRTEM with the selected area electron diffraction patterns (SAED) and histogram represents the size distribution for the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$.

Figure 4

(a-c) AFM image of the Li$_{1.1}$Co$_{0.3}$Fe$_{2.1}$O$_4$ nano ferrite sample.
Figure 5

a-b Type of sorption isotherms and pores distribution of Li1.1Co0.3Fe2.1O4.

Figure 6

a-b VSM and fitting plot between M versus 1/H2 for Li1.1Co0.3Fe2.1O4 sample.
Figure 7

SFD value for the Li1.1Co0.3Fe2.1O4 as a function of coercivity.

Figure 8
Effect of pH on the adsorption efficiency of Pb2+ ions.

Figure 9

a-b Linear fits of experimental data of the adsorption of Pb2+ onto Li1.1Co0.3Fe2.1O4 (a) Langmuir and (b) Freundlich sorption isotherms. The inset various constants of the two models.