STM-Induced SiO$_2$ Decomposition on Si(110)\(^*$

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(Received 30 January 2018; Accepted 26 July 2018; Published 11 August 2018)

Keywords: Scanning tunneling microscopy; Silicon oxides; Surface chemical reaction

Real-time scanning tunneling microscope (STM) measurements are performed during the thermal decomposition of an oxide layer on Si(110). Voids in which only oxide is removed are formed during the real-time measurements, unlike the thermal decomposition in which bulk Si is desorbed with oxide. Analysis of the STM images reveals that the measurement induces the decomposition of the oxide layer resulting from electron injection into the defect sites. The activation energy of thermal decomposition decreases by 0.4 eV in the range of 700–780°C. [DOI: 10.1380/ejssnt.2018.370]

I. INTRODUCTION

Formation and removal techniques of thin SiO$_2$ layer on Si are essential for fabrication of metal-oxide-semiconductor (MOS) field-effect transistor (FET) devices because the layer plays as gate oxide. To further develop the miniaturization and integration of MOS FET devices, thickness of the gate oxide layer must be reduced, which requires further understanding of the thermal stability of the oxide. Oxide layer decomposition also plays an important role in the fabrication of substrates for epitaxial growth because the substrates of the devices require clean surfaces that are obtained via the thermal decomposition of ultrathin oxide layers [1]. Thus, understanding the thermal stability of the oxide is crucial for the fabrication of Si-based devices.

Many studies focusing on the thermal decomposition of the oxide layer have obtained insight into the thermal stability characteristics. In particular, a time evolution analysis of morphology during the thermal decompositio of the oxide layer facilitates better understanding of the thermal decomposition mechanism [2–13]. Several evaluations revealed that a “void”, where the oxide layer is decomposed and the pure Si surface is exposed, is nucleated on the oxide layer during the initial decomposition stage [2–10]. The void is believed to grow due to SiO$_2$ decomposition at the void periphery via a reaction: SiO$_2$ + Si $\rightarrow$ 2SiO + ↑ [4–6]. Si monomers diffusing from the bottom and sidewall of the void are consumed during the reaction [7]. However, further detailed dynamics of the decomposition have not been found because these studies did not observe the decomposition in real-time.

A few real-time microscopic investigations on the morphology and dynamics of thermal decomposition of the oxide layer have been reported [11–13]. Real-time low energy electron microscope (LEEM) observations revealed that the decomposition temperature, void density and void size during the decomposition are affected by the preparation method and thickness of the oxide layer [12]. Real-time evaluation using an atomic force microscope (AFM) revealed the formation of dots in the void during the decomposition [13]. However, these studies were performed in the micrometer scale and the morphology and dynamics of thermal decomposition in the atomic-scale are not still clarified.

Real-time measurement using scanning tunneling microscope (STM) is appropriate for further understanding the time evolution of atomic-scale morphology and dynamics. For example, real-time STM study of the thermal decomposition revealed that voids are randomly nucleated on ultrathin oxide layer during annealing [11]. This indicates that the void nucleation is not depend on defects.

STM analysis is a powerful tool for obtaining atomic-scale information; however, the interpretation of resultant images can be difficult because a high electron beam dose and a high electric field are applied under the oxide layer during measurements, causing significant influence on the thermal decomposition. A previous study confirmed that an electron-beam (e-beam) of 20 eV emitted from an STM tip set at 100 nm above the surface induces selective SiO$_2$ decomposition [14] characterized by the following reaction: 2SiO$_2$ → O$_2$ + 2SiO↑ [15].

This study verifies the effect of STM-induced SiO$_2$ decomposition reaction on the thermal decomposition of the oxide layer on Si(110). Si(110) is used as the substrate because it is a candidate for high-speed MOS FET with a high hole mobility [16, 17] and has potential as a novel platform with uniformly shaped Si nanodots on it [18]. Analysis of STM images reveals that the dense e-beam emitted from the STM tip induces SiO$_2$ decomposition and consequently, the thermal decomposition process under the scanning region is affected. The STM-induced decomposition also reduces the activation energy of the SiO$_2$ decomposition.

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* This paper was presented at the 8th International Symposium on Surface Science, Tsukuba International Congress Center, Tsukuba, Japan, October 22-26, 2017.
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II. EXPERIMENTAL

A medium-doped n-type Si(110) substrate with a resistivity of 0.5–1.5 Ωcm was used. The 0.28-mm-thick wafer was cut into 1 × 7 mm² sections in the length direction, parallel to the [1\(\overline{1}2\)] direction. Surface contamination removal and oxidation were performed using Shiraki’s method [1]. This method uses hydrochloric acid to form the 5–8 Å-thick oxide layer [1]. We confirmed that the thickness of the oxide layer estimated via X-ray photoemission spectroscopy (XPS) was 5.8 Å. The thinness allowed electron tunneling, thereby implying that STM can measure the oxide surface. The sample was degassed at 550°C for 12 hr in ultrahigh vacuum (UHV) before the decomposition of the oxide layer. The oxide layer was thermally decomposed, and the morphology and electronic state evolution during the decomposition were measured by STM (JEOL JSPM-4500A) and highly sensitive synchrotron radiation XPS, respectively. Oxidized Si surfaces required a sample bias voltage > +6 V for STM measurements. The XPS experiments were conducted using SUREAC2000 at the BL23SU of Spring-8, Japan. An SR beam with a photon energy of 690 eV calibrated by measuring the Au 4f binding energy was used. The total energy resolution was < 0.25 eV, and the take-off angle of photoelectrons was 20°. These measurements were performed without taking samples out of vacuum. The sample was also annealed to nucleate voids (e.g., five days at 635°C or 28 min at 750°C). The temperature was measured using a radiation thermometer. The pressure in the UHV chamber was maintained < 3 × 10⁻⁸ Pa during the decomposition.

FIG. 1. STM images on the same scan area during the thermal decomposition of the oxide layer on Si(110) at 635°C for (a) 437-, (b) 704-, (c) 991-, and (d) 1753-min scanning durations (\(V_0 = +7.0\ V, I_1 = 0.1\ nA, 300\ nm \times 300\ nm\)). The scans were performed every 4 min. Black and white arrows denote same points, respectively. Black curves and white lines are position of shallow and deep void edges, respectively.

III. RESULTS AND DISCUSSION

Real-time STM measurement during the thermal decomposition of the oxide layer on Si(110) was performed to evaluate the STM-induced reaction’s effect on thermal decomposition. Figure 1 shows the same scan area during the thermal decomposition of oxide layer on Si(110) at 635°C. At the initial stage, several holes were formed due to silicon oxide desorption, indicating the nucleation of the void on the oxide layer [Fig. 1(a)]. The number and size of voids increased during annealing [Fig. 1(b)]. The voids were surrounded by an area that was different from the oxide surface. This area contained shallower voids than the initially formed voids because it had a striped contrast, indicating the Si(110)-16\(\times\)2 reconstructed structure. The shallower voids grew isotropically and were round, as shown in Fig. 1(c). Furthermore, several dots, such as those pointed out by arrows in Fig. 1(c), were formed in the voids. Recently, we reported that the dots are continuous Si crystal from the bulk and are formed as uniformed shape on the reconstructed structure [18]. In the final stage, all voids coalesced with each other and the entire surface exhibited the reconstructed structure [Fig. 1(d)].

STM observation revealed the nanoscale morphology and dynamics during the decomposition. However, the STM-tip affects the decomposition. The shape and depth of voids affected and not affected by STM during growth were compared to clarify the STM-tip influences.

Figure 2(a) shows an STM image of oxide Si(110) recorded halfway through decomposition at 650°C. The decomposition was performed several centimeters away from the STM tip; thus, the decomposition was not affected by STM. Voids have straight edges along the [1\(\overline{1}2\)] and [1\(\overline{1}2\)] directions, which were parallel to step rows of
void formation was the thermal decomposition of SiO$_2$ indicating the oxide layer was completely removed. The formed in the region A where was undergone different void shapes. The results indicate that the STM decomposed process near the void edge by comparing different void constructions, and the depths were 0.8 nm, 1.2 nm, and 0.4 nm than that of the reconstruction domain. The difference is also found in the void depth. Figures 2(c) and 2(d) show the height profiles corresponding to the blue lines in Figs. 2(a) and 2(b), respectively. The void, which was not affected by STM measurements, was a deep hole because the bulk Si atoms in the substrate were consumed in the reaction [Fig. 2(c)]. In contrast, the void formed during STM measurements had a shallow area [Fig. 2(d)], which was formed by the decomposition of SiO$_2$ without consuming bulk Si. Therefore, the STM-tip decomposed the oxide layer via a process different from thermal decomposition and grew voids faster.

The features of the holes in the initial stage of decomposition [white quadrangle in Fig. 2(b)] were similar to those of voids that were unaffected by STM measurements [Fig. 2(a)]. The edges were along the [112] or [1T2] directions, and the depths were ~1.5 nm. Thus, the initial void formation was the thermal decomposition of SiO$_2$ consuming bulk Si even when affected by STM. We have confirmed the difference in the decomposition process near the void edge by comparing different void shapes. The results indicate that the STM decomposed the oxide layer. The change in the height of the oxide surface versus the number of STM scans was analyzed to clarify influence of decomposition by the STM on oxide layer thickness.

Figure 3(a) is STM image which displays three regions undergone different numbers of STM scans, as labeled with A/B/C. The 16 × 2 reconstructed structure was formed in the region A where was undergone > 300 scans indicating the oxide layer was completely removed. The blue line transverses the region B (25 scans) and C (5 scans) and indicates the position of the height profile, as shown in Fig. 3(b). The height decreased as the number of scans increased, indicating that the STM measurement etched the oxide layer. This suggests that the STM decomposes the oxide layer and make it thin. This finally completely removes oxide resulting that pure Si surface is exposed such as the region A. In addition, the edge of oxide layer is easier to decompose than the terrace because there is always next to defect. This makes shallow part around deep void.

A possible cause of the decomposition of the oxide layer is the tunnel current from the STM-tip. A previous study reported that an e-beam of 20 eV decomposed the thermal oxide layer [14]. The SiO$_2$ decomposition by the e-beam is believed to follow the reaction 2SiO$_2$ → O$_2$ + +2SiO ↑ [15] and does not consume bulk Si atoms to reduce SiO$_2$. Furthermore, the extreme condition of STM observation (bias voltage $V_s$ = +10 V, tunnel current $I_t$ = 6 nA) created a leakage site of current on the thermal oxide layer [20]. For this phenomenon to occur, it is believed that the injection of electrons into a defect, which acts as an electron trap site, is essential [20]. The bond energy of Si–O in SiO$_2$ is 6.5 eV; therefore, the current injection from the STM-tip can cut the Si–O bond. However, the STM could not decompose the thermally oxidized surface [20]. The difference should be caused by the difference of defect density. The oxide used in this study was formed via Shiraki’s method. This forms oxide layer containing many defects. The defects make the oxide layer easier to decomposed than thermal oxide layer [1, 21]. Furthermore, the sample was annealed for a long time period before STM measurement to nucleate voids. This should also make starting point of decomposition by the STM.

Although the electric field between the surface and the STM-tip was very high, field evaporation would not cause decomposition of oxide layer. The electric field between the surface and the tip was ~7 × 10$^9$ V/m because the distance and the sample bias voltage were ~1 nm and $V_s$ = 7 V, respectively. The electric field was almost identical to the condition of typical W tip field evaporation, which was reported to be 4 × 10$^{10}$ V/m via field ion microscope [19]. However, since the electric field does not affect the decomposition of the oxide layer [20], the electric field is not the mechanism of decomposition.

The STM-tip did not mechanically etch the oxide surface. Although the etching proceeded in the scan direction, the void grew isotropically. Furthermore, mechanical etching would deposit shavings at the end of scan but such a deposition was not found in Fig. 3(a).

The activation energy of oxide layer decomposition affected by STM measurement was compared with that of purely thermal decomposition of oxide layer estimated using XPS to quantitatively evaluate the effect of STM. The analysis was performed at a temperature > 700°C. The reconstructed structure cannot exist at this temperature because phase transition to 11 structure occurs at 695°C [22]. Thus, obtained activation energy is not affected by formation of reconstructed structure.

The activation energy with less influence other than thermal decomposition was estimated via XPS because the electrons traveling from the oxide to the cathode cannot form a leakage site [20]. The activation energy was
Figure 4. (a) The time evolution of residual oxide during oxide layer decomposition on Si(110). (b) The Arrhenius plot of thermal decomposition rates estimated at 750°C, 760°C, and 780°C as a function of 10000/T.

Figure 5. Arrhenius plot of void growth rates estimated at 700°C, 735°C, and 740°C as a function of 10000/T.

determined from the reduction rates of oxygen estimated from O 1s spectra measured via XPS. The amount of oxygen remained constant (for example, 28 min at 750°C) after annealing began. Figure 4(a) shows the residuals of oxygen with respect to the initial amount. The beginning of the changing oxygen amount indicated the nucleation of a void. Since the number of voids increased, the reduction rate of the residual accelerated during the first half of the decomposition. However, in the second half of the decomposition, the reduction rate decelerated because the void edge wherein the decomposition reaction occurred was shortened by combining voids. Therefore, the maximum rate, where the acceleration and the deceleration were balanced, was adopted to determine the activation energy.

Figure 4(b) shows the Arrhenius plot of thermal decomposition rates estimated at 750°C, 760°C, and 780°C as a function of 10000/T. The solid line indicates the result of fitting the data to the Arrhenius form, \( R = R_0 \exp(-E_a/k_B T) \), where \( R \) is the decomposition rate and \( E_a \) is the activation energy. The activation energy obtained from the plots was 3.0 eV.

This activation energy is 1.2 eV lower than that of the thermal oxide layer on Si(100) determined via LEEM observation [21] possibly because of the oxidation method and surface orientation of the Si substrate. The oxide formed using hydrochloric acid used herein included several defects derived from Cl inclusion, leading to low temperature decomposition compared to thermal oxide [1, 21]. The Si(110) surface was rapidly oxidized in an oxygen atmosphere, indicating the low activation energy of the Si + O → SiO reaction on the Si(110) surface [23, 24]. Since the free energy of Si, O₂, and SiO₂ do not depend on surface orientation, the activation energy of decomposition should also be low.

Figure 5 shows the Arrhenius plot of void growth rate evaluated from STM observation at 700°C, 725°C, and 740°C as a function of 10000/T with a +6 V sample bias voltage. The solid line indicates the result of fitting the date to the Arrhenius form, \( r = r_0 \exp(-E_a/k_B T) \), where \( r \) is the differentiation of void diameter. The activation energy obtained from the plots was 2.6 eV, which was 0.4 eV lower than that determined via XPS measurements. Therefore, STM measurements reduce the activation energy of oxide layer decomposition on Si(110) by 0.4 eV.

IV. CONCLUSIONS

The influence of STM measurements on the decomposition of the oxide layer on Si(110) was analyzed. The STM tip etched the oxide surface and affected on the void shape. The decomposition reaction induced via STM should occur at defects, where the defect site acts as an electron trap and the Si–O bonds were weak. From comparison to XPS estimation, the resultant activation energy reduction of the decomposition was estimated to be 0.4 eV upon the influence of STM measurements.

ACKNOWLEDGMENTS

The XPS experiments were performed at the BL23SU of Spring-8 with the approval of the Japan Synchrotron Radiation Research Institute (JASRI) (Proposal Nos. 2016B3801 and 2016A3801).

[1] A. Ishizaka and Y. Shiraki, J. Electrochem. Soc. 133, 666 (1986).

[2] D. Jones and V. Palermo, Appl. Phys. Lett. 80, 673 (2002).
[3] K. E. Johnson, P. K. Wu, M. Sander, and T. Engel, Surf. Sci. 290, 213 (1993).
[4] H. Watanabe, K. Fujita, and M. Ichikawa, Appl. Phys. Lett. 70, 1095 (1997).
[5] K. E. Johnson and T. Engel, Phys. Rev. Lett. 69, 339 (1992).
[6] G. W. Rubloff, J. Vac. Sci. Technol. A 8, 1857 (1990).
[7] R. Tromp, G. W. Rubloff, P. Balk, F. K. LeGoues, and E. J. van Loenen, Phys. Rev. Lett. 55, 2332 (1985).
[8] K. Fujita, H. Watanabe, and M. Ichikawa, Appl. Phys. Lett. 70, 2807 (1997).
[9] N. Miyata, H. Watanabe, and M. Ichikawa, Phys. Rev. Lett. 84, 1043 (2000).
[10] K. Fujita, H. Watanabe, and M. Ichikawa, Surf. Sci. 398, 134 (1998).
[11] Y. Wei, R. M. Wallace, and A. C. Seabaugh, Appl. Phys. Lett. 69, 1270 (1996).
[12] H. Hibino, M. Uematsu, and Y. Watanabe, J. Appl. Phys. 100, 113519 (2006).
[13] D. Fujita, J. Surf. Sci. Soc. Jpn. 36, 459 (2015) (in Japanese).
[14] H. Iwasaki, T. Yoshinobu, and K. Sudo, Nanotechnology 14, R55 (2003).
[15] S. Fujita, S. Maruno, H. Watanabe, and M. Ichikawa, Appl. Phys. Lett. 69, 638 (1996).
[16] T. Satô, Y. Takeishi, H. Hara, and Y. Okamoto, Phys. Rev. B 4, 1950 (1971).
[17] W. Cheng, A. Teramoto, M. Hirayama, S. Sugawa, and T. Ohmi, Jpn. J. Appl. Phys. 45, 3110 (2006).
[18] M. Yano, Y. Uozumi, S. Yasuda, and H. Assoka, Jpn. J. Appl. Phys. 57, 06HD04 (2018).
[19] H.-W. Fink, IBM J. Res. Dev. 30, 460, (1986).
[20] H. Watanabe, K. Fujita, and M. Ichikawa, Appl. Phys. Lett. 72, 1987 (1998).
[21] C. Detavernier, R. L. Van Meirhaeghe, F. Cardon, and K. Maex, Thin Solid Films 386, 19 (2001).
[22] Y. Yamamoto, T. Sueyoshi, T. Sato, and M. Iwatsuki, Surf. Sci. 466, 183 (2000).
[23] H. Togashi, Y. Takahashi, A. Kato, A. Konno, H. Asaoka, and M. Suemitsu, Jpn. J. Appl. Phys. 46, 3239 (2007).
[24] M. Suemitsu, A. Kato, H. Togashi, A. Konno, Y. Yamamoto, Y. Teraoka, A. Yoshigoe, Y. Narita, and Y. Enta, Jpn. J. Appl. Phys. 46, 1888 (2007).