LIMIT CYCLES FOR PLANAR SEMI-QUASI-HOMOGENEOUS POLYNOMIAL VECTOR FIELDS

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Abstract. This paper is concerned with the limit cycles for planar semi-quasi-homogeneous polynomial systems. We give some explicit criteria for the nonexistence and existence of periodic orbits. Let \( N = N(p, q, m, n) \) be the maximum number of limit cycles of such system. A lower bound is given for \( N = N(p, q, m, n) \). The cyclicity and center problem are studied for some subfamilies of semi-quasi-homogeneous polynomial systems. Our results generalize those obtained for polynomial semi-homogeneous systems.

1. Introduction and statement of the main results

A function \( f(x, y) \) is called a \((p, q)\)-quasi-homogeneous function of weighted degree \( m \) if \( f(\lambda^p x, \lambda^q y) = \lambda^m f(x, y) \) for all \( \lambda \in \mathbb{R} \). Let \( P_m(x, y) \) and \( Q_n(x, y) \) be \((p, q)\)-quasi-homogeneous polynomials of weighted degree \( p - 1 + m \) and \( q - 1 + n \), respectively. We say that \( X = (P_m(x, y), Q_n(x, y)) \) is a planar \((p, q)\) semi-quasi-homogeneous polynomial vector field if \( m \neq n \). The system of differential equations associated to \( X \) is

\[
\begin{align*}
\frac{dx}{dt} &= P_m(x, y) = \sum_{p_i + q_j = p + m - 1} a_{ij} x^i y^j, \\
\frac{dy}{dt} &= Q_n(x, y) = \sum_{p_i + q_j = q + n - 1} b_{ij} x^i y^j.
\end{align*}
\]

Here \( p, q, m \) are positive integers and \( P_m(x, y) \) and \( Q_n(x, y) \) are coprime in the ring \( \mathbb{R}[x, y] \). To be short we denote this by \((P(x, y), Q(x, y)) = 1\).

Observed that the above definition is the natural one for the following reason:

(i) When \( p = q = 1 \), \( X \) is called a semi-homogeneous vector field\(^3\) \(^4\). That is to say, the above definition coincides with the usual definition of the homogeneous cases.

(ii) If \( m = n \), then \( X \) is called a \((p, q)\)-quasi-homogeneous vector field of weighted degree \( m \), see \(^1\), Chapter 7. In particular if \( p = q = 1 \), \( X \) is a homogeneous polynomial vector field of degree \( m \).

(iii) The finite (resp. infinity) singular points of semi-homogeneous vector field can be studied by homogeneous blow-up (resp. Poincaré compactification)

2000 Mathematics Subject Classification. Primary 34C05, 34A34, 34C14.
Key words and phrases. Limit cycles, semi-quasi-homogeneous vector fields.
Supported by the NSF of China (No.11171355), the Ph.D. Programs Foundation of Ministry of Education of China (No. 20100171110040) and Program for New Century Excellent Talents in University.
whereas the one of semi-quasi-homogeneous vector fields can be studied by the use of quasi-homogeneous blow-up (resp. Poincaré-Lyapunov compactification)\cite{10}.

One of the classical problem in the qualitative theory of planar polynomial systems is the study of the number of limit cycles, which is known as the 16th Hilbert’s problem. There has been a substantial amount of work devoted to solving this problem for homogeneous, quasi-homogeneous or semi-homogeneous polynomial vector fields. W. Li etc.\cite{13} provided an upper bound for the maximum number of limit cycles bifurcating from the period annulus of any homogeneous and quasi-homogeneous centers, which can be obtained using the Abelian integral method. Gavrilov etc.\cite{11} gave a more general results on the limit cycles bifurcated from the periodic orbits of quasi-homogeneous centers. Cima and Llibre\cite{7} investigated the algebraic and topological classification of homogeneous cubic vector fields. In 1997, Cima, Gasull and Manosas\cite{6} studied the limit cycles for semi-homogeneous vector fields. They proved that semi-homogeneous systems can exhibit periodic orbits only when \(nm\) is odd and there exists such system with at least \((n+m)/2\) limit cycles. In the papers\cite{3,4} Cairó and Llibre study the phase portraits of semi-homogeneous vector fields with \((m,n) = (1,2)\) and \((m,n) = (1,3)\) respectively. Some papers are concerned with the integrability\cite{9,16} and structural stability\cite{14,15,17} for quasi-homogeneous or semi-quasi-homogeneous polynomial vector fields.

It is always be possible to decompose a vector field \(\mathbf{X}\) in a formal series centered in \(0 \in \mathbb{R}^n\): \(\mathbf{X} = \sum_{j \geq k} X_j\), where \(X_j\) are \((p,q)\)-quasi-homogeneous vector field of weighted degree \(k\), \(k\) is the first integer such that \(X_k \neq 0\)\cite{2}. Therefore in order to get more information on limit cycles of general vector fields \(\mathbf{X}\) it is essential to study the geometrical properties of quasi-homogeneous and semi-quasi-homogeneous vector field. We note that Coll, Gasull and Prohens have investigated planar vector fields defined by the sum of two quasi-homogeneous vector fields\cite{8}.

In this paper we study the limit cycles of planar semi-quasi-homogeneous polynomial vector fields, defined in \(\mathbf{1}\). To simplify the statements of the main results, first of all we give the following lemma, which will be proved in section\cite{2}

\textbf{Lemma 1.} Suppose \((p,q) = k \geq 2\) in \(\mathbf{1}\), then there exists a unique vector \((p',q',m',n')\) with \((p',q') = 1\) such that system \(\mathbf{1}\) is a \((p',q')\) semi-quasi-homogeneous vector field.

By Lemma\cite{1} we suppose that the following conventions hold without loss of generality.

\textbf{Convention} We always assume in this paper that

(i) \(p\) is odd and \((p,q) = 1\), and

(ii) \(P_m(x,y)\) and \(Q_m(x,y)\) are coprime polynomials.

Our main results are the following four theorems. The first one provides explicit criteria for the nonexistence and existence of periodic orbits, and a lower bound of \(N = N(p,q,m,n)\) for the maximum number of limit cycles of system \(\mathbf{1}\).
Theorem 2. Let $P_m(x, y)$ and $Q_n(x, y)$ are $(p, q)$-quasi-homogeneous polynomials of degree $p - 1 + m$ and $q - 1 + n$, respectively, $m \neq n$. For system (1), the following statements hold:

(i) If either $q \nmid p + m - 1$ or $p \nmid q + n - 1$, then system (1) has no periodic orbit.

(ii) If $q|p + m - 1$ and $p|q + n - 1$, then system (1) is reduced to the following normal form

\[
\frac{dx}{dt} = P_m(x, y) = \sum_{i=0}^{[r_1/p]} a_i x^{iq} y^{r_1-ip}, \quad \frac{dy}{dt} = Q_n(x, y) = \sum_{j=0}^{[r_2/q]} b_j x^{r_2-jq} y^{jp}.
\]

where $r_1 = (p + m - 1)/q$, $r_2 = (q + n - 1)/p$, $[r_1/p]$ denotes the integer part of $r_1/p$.

If system (2) has periodic orbit, then one of the following conditions holds:

(ii.1) If both $p$ and $q$ are odd, then $m$, $n$, $r_1$ and $r_2$ are odd.

(ii.2) If $p$ is odd and $q$ is even, then $m$ and $n$ are even, $r_1$ and $r_2$ are odd, respectively.

(iii) If (ii.1) holds, then there exist $P_m(x, y)$ and $Q_n(x, y)$ such that each of the following assertions holds:

(a) All solutions of (2) are periodic.

(b) No solutions of (2) are periodic.

(c) There are periodic and non-periodic solutions of system (2).

Let $N = N(p, q, m, n)$ be the maximum number of limit cycles of system (2). Then $N \geq \left[\left(\left\lfloor \frac{r_1}{p} \right\rfloor + 1 \right)/2 \right] + \left[\left\lfloor \frac{r_2}{q} \right\rfloor + 1 \right)/2 - 1$.

(iv) If (ii.2) holds and system (2) has periodic orbit, then the origin is a center.

The results in Theorem 2 generalize those obtained for semi-homogeneous polynomial vector fields, proved in [6].

One of important problem in the qualitative theory of planar polynomial vector fields is the study of the local phase portrait at the singularities to characterize when a singular is a center, which is called the center problem. A singular point $O$ of a real planar analytic vector field $X$ is called a center if there exists a neighborhood $U$ of $O$ such that $U \setminus \{O\}$ is filled with periodic integral curves of the vector field. The center is called elementary if the linearization $dX(O)$ of the vector field is a rotation, otherwise the center is called non-elementary. If the origin is a center of system (2), then it is in general non-elementary. The cyclicity is the total number of limit cycles which can emerge from a configuration of trajectories (center point, period annulus, separatrix cycle) under a perturbation.

The center and its cyclicity problem are only solved for some special kind of planar polynomial systems: quadratic systems, cubic systems with homogeneous nonlinearities, and so on. In the following two theorems, we give some more information on the center and cyclicity problem for some subfamily of the semi-quasi-homogeneous system (2).
Theorem 3. Suppose that \( q | p + m - 1, p | q + n - 1 \) and the condition (ii.1) in Theorem 2 holds. Consider the system

\[
\frac{dx}{dt} = \sum_{i=0}^{[r_1/p]} a_i x^{i} y^{r_1-1-\tilde{p}}, \quad \frac{dy}{dt} = b_0 x^{r_2}
\]

with \( r_1 > r_2 \) and \( m > n \). If \( a_0 b_0 < 0 \), then the following statements hold.

(i) The origin is a center if and only if \( a_1 = a_3 = \cdots = a_\kappa = 0 \), where

\[
\kappa = \begin{cases} 
\left[ \frac{r_1}{p} \right], & \text{if } \left[ \frac{r_1}{p} \right] \text{ is odd,} \\
\left[ \frac{r_1}{p} \right] - 1, & \text{if } \left[ \frac{r_1}{p} \right] \text{ is even.}
\end{cases}
\]

(ii) If \( a_1 = a_3 = \cdots = a_{2i-3} = 0, a_{2i-1} \neq 0, 2i - 1 \leq \kappa \), then the origin is a weak focus and under perturbations inside the family it can produce \( i - 1 \) limit cycles.

(iii) The cyclicity of the center is \( (\kappa - 1)/2 = \lfloor (r_1/p) - 1 \rfloor/2 \).

In the next theorem we study the local phase portraits at the origin for system (2) with \( m = 1 \).

Theorem 4. Suppose that \( m = 1 \) in (2) and the condition (ii.1) in Theorem 2 holds. Then system (2) is reduced to

\[
\frac{dx}{dt} = P_1(x, y) = a_1 x + a_0 y^p, \quad \frac{dy}{dt} = Q_n(x, y) = \sum_{j=0}^{n/p} b_j x^{n/p-1} y^p
\]

where \( n/p \) is an odd number and \( (P_1(x, y), Q_n(x, y)) = 1 \).

(i) Suppose that \( a_0 = 0, a_1 \neq 0 \), then \( b_{n/p} \neq 0 \). Moreover the origin is an unstable (resp. stable) node if \( a_1 > 0, b_{n/p} > 0 \) (resp. \( a_1 < 0, b_{n/p} < 0 \) ) and a saddle if \( a_1 b_{n/p} < 0 \) respectively.

(ii) Suppose \( a_1 \neq 0, a_0 = 0 \).

(a) If \( p = 1 \), then

(a.1) If \( \sum_{j=0}^{n} (-1)^{n-j} b_j/(a_0^{j-1} a_1^{n-j+1}) > 0 \) and \( a_1 > 0 \) (resp. \( a_1 < 0 \)), then the origin is a unstable (resp. stable) node.

(b.1) If \( a_1 (\sum_{j=0}^{n/p} b_j (-a_0/a_1)^{n/p-j}) > 0 \) and \( a_1 > 0 \) (resp. \( a_1 < 0 \)), then the origin is a unstable (resp. stable) node.

(b) If \( p \geq 3 \), then

(b.1) If \( a_1 (\sum_{j=0}^{n/p} b_j (-a_0/a_1)^{n/p-j}) > 0 \) and \( a_1 > 0 \) (resp. \( a_1 < 0 \)), then the origin is a unstable (resp. stable) node.

(b.2) If \( a_1 (\sum_{j=0}^{n/p} b_j (-a_0/a_1)^{n/p-j}) < 0 \), then the origin is a saddle.

(iii) Suppose \( a_0 b_0 < 0, a_1 = 0, n > p^2 \), then the origin is a center if and only if \( b_1 = b_3 = \cdots = b_{n/p} = 0 \).

Finally, we study the cyclicity at infinity for a kind of dual version of the above vector field. Before that we introduce the \((l_1, l_2)\)-trigonometric function which were given by Lyapunov [12] in his study of the stability of degenerate singular
points. Let \( z(\phi) = C\phi \) and \( \omega(\phi) = S\phi \) be the solution of the following initial problem

\[
\dot{z} = -\omega^{2l_1 - 1}, \quad \dot{\omega} = z^{2l_2 - 1}, \quad z(0) = \frac{1}{l_1}, \quad \omega(0) = 0.
\]

For \((l_1, l_2) = (1, 1)\), we have that \( C\phi = \cos \phi, \) \( S\phi = \sin \phi \), i.e. the \((1, 1)\)-trigonometric functions are the classical ones.

Consider the system

\[
\begin{align*}
\frac{dx}{dt} &= -y^{r_1}, \\
\frac{dy}{dt} &= x^{r_2} + \sum_{j=1}^{[r_2/q]} b_j x^{r_2 - jq} y^{jp},
\end{align*}
\]

where \( r_1 = 2l_1 - 1 \) and \( r_2 = 2l_2 - 1 \) are defined in Theorem 2.

In order to study the behavior of the trajectories of a planar differential system near infinity it is possible to use a compactification. In this paper we prefer to reduce the "infinity" to a single point. The construction of this compactification is as follows (see for instance [5]): we first write system (5) in the generalized polar coordinate \((\rho, \phi)\) by the transformation

\[
\begin{align*}
x &= \rho^{l_1} C\phi, \\
y &= \rho^{l_2} S\phi,
\end{align*}
\]

then change \( \rho \) by \( 1/\rho \) and finally reparametrize the system to obtain a smooth one.

It is said that a planar differential equation has a center at infinity if there is a compact set of \( \mathbb{R}^2 \) such that outside it all trajectories of that differential equations are closed [5].

**Theorem 5.** Suppose that \( q|p + m - 1, \) \( p|q + n - 1, \) \( r_1 > r_2, \) \( m > n \) and the condition (ii.1) in Theorem 2 holds. For system (5), the following statements hold.

(i) A neighborhood of infinity is filled with periodic orbits if and only if \( b_1 = b_3 = \cdots = b_\mu = 0, \) where

\[
\mu = \begin{cases} 
\left[ \frac{r_2}{q} \right], & \text{if } \left[ \frac{r_2}{q} \right] \text{ is odd,} \\
\left[ \frac{r_2}{q} \right] - 1, & \text{if } \left[ \frac{r_2}{q} \right] \text{ is even.}
\end{cases}
\]

(ii) If \( b_1 = b_3 = \cdots = b_{2i-3} = 0, \) \( b_{2i-1} \neq 0, \) \( 2i - 1 \leq \mu, \) then the stability of the infinity is given by the sign of \( b_{2i-1} \). Moreover the maximum number number of periodic orbits that bifurcate from infinity inside the family for small perturbations is \( i - 1 \).

(iii) The cyclicity of the infinity is \((\mu - 1)/2 = [([r_2/q] - 1)/2]\).

Theorem 3 and Theorem 5 generalize those for semi-homogeneous polynomial vector fields in [6].
2. SOME PRELIMINARY RESULTS

In this section we give some preliminary results which will be used in the proof of the main results. First of all we prove Lemma 1.

Proof of Lemma 1. Let $p = kp', q = kq'$. Since $(p, q) = k \geq 2$, we have $(p', q') = 1$. It follows from the definition of $P_m(x, y)$ in (1) that

$$kp' - 1 + m = kp'i + kq'j,$$

which implies that $k|m - 1$. Using the same arguments one obtains $k|n - 1$. Taking $m' = 1 + (m - 1)/k$ and $n' = 1 + (n - 1)/k$, the assertion follows. □

Lemma 6. [6, 12] $C\phi$ and $S\phi$, defined in Section 1, have the following properties.

(i) $l_1 C_{2l_2} \phi + l_2 S_{2l_1} \phi = 1$.

(ii) $dS\phi/d\phi = C_{2l_2} - 1 \phi$, $dC\phi/d\phi = -S_{2l_1} - 1 \phi$.

(iii) $C(-\phi) = C\phi$, $S(-\phi) = -S\phi$.

(iv) $C\phi$ and $S\phi$ are $T$-periodic functions with

$$T = T_{1, l_2} = 2l_1^{-\frac{1}{2l_2}} l_2^{-\frac{\pi_1}{2l_1}} \frac{\Gamma(\frac{1}{2l_1}) \Gamma(\frac{1}{2l_2})}{\Gamma(\frac{1}{2l_1} + \frac{1}{2l_2})},$$

where $\Gamma$ denotes the gamma function.

Lemma 7. [6] $\int_0^T S_{\alpha} \phi C_{\beta} \phi d\phi \neq 0$ if and only if $\alpha$ and $\beta$ are even.

Finally, we give two theorems which characterize the local phase portraits at the singular points.

Theorem 8. [10] Let $(0, 0)$ be an isolated singular point of the vector field $X$ given by

$$\dot{x} = A(x, y), \quad \dot{y} = \lambda y + B(x, y),$$

where $A(x, y)$ and $B(x, y)$ are analytic in a neighborhood of the origin with $A(0, 0) = B(0, 0) = DA(0, 0) = DB(0, 0) = 0$ and $\lambda > 0$. Let $y = f(x)$ be the solution of the equation

$$\lambda y + B(x, y) = 0$$

in a neighborhood of the origin, and suppose the function $A(x, f(x))$ has the expression

$$A(x, f(x)) = a_m x^m + o(x^m),$$

where $m \geq 2$ and $a_m \neq 0$. Then there always exists an invariant analytic curve, called the strong unstable manifold, tangent at $(0, 0)$ to the $y$-axis, on which $X$ is analytically conjugate to $\dot{y} = \lambda y$; it represents repelling behavior since $\lambda > 0$. Moreover the following statements hold:

(i) If $m$ is odd and $a_m < 0$, then $(0, 0)$ is a topological saddle.

(ii) If $m$ is odd and $a_m > 0$, then $(0, 0)$ is an unstable topological node.

(iii) If $m$ is even, then $(0, 0)$ is a saddle-node, that is, a singular point whose neighborhood is the union of one parabolic and two hyperbolic sectors.
We say that a family of vector fields \( X_\lambda, \lambda \in \mathbb{R}^k \) is **stable by homoteces** if all \( \lambda \in \mathbb{R}^k \) and for all \( \epsilon \in \mathbb{R} \) there exists \( \lambda' \in \mathbb{R}^k \) such that \( \epsilon X_\lambda = X_{\lambda'} \). Fixed \( p, q \), any analytic vector field \( X \), such that \( X(0) = 0 \), can be decomposed as

\[
X = \sum_{i=1}^{\infty} h^{(p,q)}_i(X),
\]

where \( h^{(p,q)}_i(X) \) is \((p,q)\)-quasi-homogeneous vector fields of weighted degree \( i \).

**Theorem 9.** [5] Consider the vector field

\[
X = -y^n \frac{\partial}{\partial x} + x^m \frac{\partial}{\partial y} + Y(x, y, c) + \sum_{i=1}^{l} b_i Z_i(x, y),
\]

where \( n = 2p - 1 \geq m = 2q - 1 \), are natural numbers, \( c \in \mathbb{R}^u \), \( u \in \mathbb{N} \), and \( b = (b_1, b_2, \cdots, b_l) \in \mathbb{R}^l \), \( l \in \mathbb{N} \), and

(i) \( Y(x, y, c) = Y_1(x, y, c)(\partial/\partial x) + Y_2(x, y, c)(\partial/\partial y) \) has analytic components, \( Y(0, 0, c) = 0 \), \( h^{(p,q)}_i(Y(x, y, c)) = 0 \) for all \( i \leq k = 2pq - p - q + 1 \) (the weighted degree of \(-y^n(\partial/\partial x) + x^m(\partial/\partial y)\)) and the differential equations associated \( b = 0 \) has a center at the origin.

(ii) \( Z_i \) are \((p,q)\)-quasi-homogeneous vector fields of weighted degree \( d_i \), greater than \( k = 2pq - p - q + 1 \) and \( d_i > d_j \) when \( i > j \). Furthermore

\[
I_i = \iint_{px^{2q} + qy^2 \leq 1} \text{div} Z_i(x, y) dxdy \neq 0
\]

for all \( i \).

Then for the differential equations associated to (8) the following hold.

(a) The origin is a center if and only if \( b = 0 \).

(b) If \( b_1 = b_2 = \cdots = b_{l-1} = 0 \) and \( b_i \neq 0 \) the stability of the origin is given by the sign \( b_i I_i \) and its cyclicity inside the family (8) is \( i - 1 \).

(c) The cyclicity of the center inside the family (8) is \( l - 1 \).

(d) If in addition the family \( Y(x, y, c) \) is stable by homoteces, then the maximum number of limit cycles that can bifurcate for \( \epsilon \) small enough, from the curve \( px^{2q} + qy^{2p} = h \) for the vector field

\[
-y^n \frac{\partial}{\partial x} + x^m \frac{\partial}{\partial y} + \epsilon \left( Y(x, y, c) + \sum_{i=1}^{l} b_i Z_i(x, y) \right),
\]

is \( l - 1 \). Moreover there exist \( b \in \mathbb{R}^l \) and \( c \in \mathbb{R}^u \) such that this number is attained.

The following remark was given in [5].

**Remark 10.** If we change the hypotheses of Theorem 9 by

(i') \( Y(x, y, c) \) has polynomial components, \( h^{(p,q)}_i(Y(x, y, c)) = 0 \) for all \( i \geq k = 2pq - p - q + 1 \) and the differential equations associated \( b = 0 \) has a center at infinity.
(ii') \( Z_i \) are \((p,q)\)-quasi-homogeneous vector fields of weighted degree \( d_i, d_i < d_j \) when \( i > j \), and \( d_i < k = 2pq - p - q + 1 \), and furthermore
\[
I_i = \iint_{p^2 + qy^2 \leq 1} \text{div} Z_i(x,y) dxdy \neq 0 \quad \text{for all } i,
\]
then the same conclusions hold, changing the word "origin" to "infinity".

The above two theorems and remark are important for the proof of Theorem 3.5.

3. Proof of Theorem 2

In this section we will prove Theorem 2.

Lemma 11. The origin is a unique (real and finite) singular point of system (1).

Proof. By definition of \((p,q)\) semi-quasi-homogeneous systems, we have that
\[
P_m(\lambda^p x, \lambda^q y) = \lambda^{p+m-1} P_m(x,y), \quad Q_n(\lambda^p x, \lambda^q y) = \lambda^{q+n-1} Q_n(x,y).
\]
Taking \( \lambda = 0 \) into (9) for fixed \((x,y)\), one obtains \( P_m(0,0) = Q_n(0,0) = 0 \). Therefore the origin is a singular point of system (1).

Let \((a,b)\) be the singular point of system (1) with \((a,b) \neq (0,0)\). It follows from (9) that \( P_m(\lambda^p a, \lambda^q b) = Q_n(\lambda^p a, \lambda^q b) = 0 \). This yields that the curve \(\{(\lambda^p a, \lambda^q b) : \lambda \in \mathbb{R}\}\) is fulfill of singular points, in contradiction that \( P_m(x,y) \) and \( Q_n(x,y) \) are coprime. \(\square\)

Proof of Theorem 2(i). Suppose that \( P_m(x,y) \) is defined in (1). The pair \((i,j)\) in the expression of \( P_m(x,y) \) is a non-negative solution of the equation
\[
pi + qj = p - 1 + m.
\]
As usual, we use the notation non-negative integer (resp. integer) solution if \( i \) and \( j \) are non-negative integers (resp. integers). If \( q \nmid p + m - 1 \), then there does not exist \( j \) such that \((0,j)\) satisfies (10). Therefore \( P_m(0,0) = 0 \), which implies that \( x = 0 \) is an invariant lines of system (1). Since Lemma 11 shows that the origin is a unique singular point of system (1), system (1) has no periodic orbit.

If \( p \nmid q + n - 1 \), then by the same arguments as above we conclude that \( y = 0 \) is an invariant lines of system (1). Therefore system (1) has no periodic orbit. \(\square\)

To prove the assertion (ii) of Theorem 2 we give the following lemma.

Lemma 12. Suppose that the Convention (i) holds, \( q \) is even and \( n \) is odd, then system (2) has no periodic orbit.

Proof. By direct computation we have
\[
\left. \frac{dy}{dt} \right|_{y=0} = Q_n(x,0) = b_0 x^{r2} = b_0 (x^{1/p})^{q-1+n}.
\]
If \( b_0 \neq 0 \) and \( n \) is odd, then \( (dy/dt)|_{y=0} \) does not change the sign. Since the flow on \( x\)-axis have the same direction in the phase plane, there is no periodic orbit surrounding the origin.
If \( b_0 = 0 \), then \( y = 0 \) is an invariant line of system (2), which yields the assertion. This finishes the proof. \( \square \)

Proof of Theorem 2(ii). If \( q \mid p + m - 1 \), then \((0, r_1)\) is a non-negative integer solution of equation (10). Moreover, all the integer solution of equation (10) are \( \{(iq, r_1 - ip) : i \in \mathbb{Z}\} \), which implies that

\[
P_m(x, y) = \sum_{i=0}^{[r_1/p]} a_ix^{iq}y^{r_1-ip}.
\]

By the same arguments, we have \( Q_n(x, y) = \sum_{j=0}^{[r_2/q]} b_jx^{r_2-jq}y^{jp} \).

Suppose that system (2) has at least one periodic orbit. Note that system (2) has a unique singular point at the origin and hence all periodic orbits surround the origin.

We split the proof into two cases.

(a) Assume that \( p \) and \( q \) are odd. By applying the definition of a \((p, q)\) semi-quasi-homogeneous system, we have

\[
P_m(-x, -y) = P_m((-1)^px, (-1)^qy) = (-1)^{p-1+m}P_m(x, y) = (-1)^mP(x, y).
\]

If \( m \) is even, then \( P_m(-x, -y) = P(x, y) \), which implies that \( dx/dt \mid_{x=0} = P(0, y) = P(0, -y) \). Hence the flow on \( x = 0 \) goes either from left to right, or from right to left in the phase plane, in contradiction with the assumption that system (2) has at least a period orbit. Therefore \( m \) is odd, which gives \( P_m(-x, -y) = -P_m(x, y) \).

It follows from (2) that

\[
P_m(-x, -y) = \sum_{i=0}^{[r_1/p]} a_i(-1)^{r_1+(q-p)i}x^{iq}y^{r_1-ip} = -\sum_{i=0}^{[r_1/p]} a_ix^{iq}y^{r_1-ip} = -P_m(x, y).
\]

Since both \( p \) and \( q \) are odd, one obtains \( a_i((-1)^{r_1} + 1) = 0 \). If \( r_1 \) is even, then \( a_i = 0 \) for \( i = 0, 1, 2, \cdots, [r_1/p] \), i.e. \( P_m(x, y) \equiv 0 \). In this case system (2) has no periodic orbit. Therefore \( r_1 \) is odd.

Using the same arguments as above, we conclude that both \( n \) and \( r_2 \) are odd. The statement (ii.1) is proved.

(b) Assume that \( p \) is odd and \( q \) is even. By the definition of \( r_1 \) and \( r_2 \), we conclude that \( m \) is even, and \( r_2 \) is odd (resp. even) if and only if \( n \) is even (resp. odd). Since we suppose that system (2) has at least one periodic orbit, it follows from Lemma 12 that \( n \) is even, which implies that \( r_2 \) is odd. In what follows we only need to prove that \( r_1 \) is odd.

Note that

\[
\frac{dx}{dt} \bigg|_{x=0} = a_0y^{r_1}.
\]

If \( r_1 \) is even, then \( (dx/dt) \big|_{x=0} \) does not change the sign or equals to zero identically, which implies that system (2) has no periodic orbit. The assertion (ii.2) follows. \( \square \)

The following theorem will be used to prove the assertion (iii) of Theorem 2.
**Theorem 13.** Suppose that (ii.1) of Theorem 2 holds and \( q|p+m-1, \ p|q+n-1 \), then there exist \( P_m(x, y) \) and \( Q_n(x, y) \) such that system

\[
\frac{dx}{dt} = y^{r_1} + \epsilon P_m(x, y), \quad \frac{dy}{dt} = -x^{r_2} + \epsilon Q_n(x, y)
\]

has at least \( [(r_1/p) + 1]/2 + [(r_2/q) + 1]/2 - 1 \) limit cycles, where \( \epsilon \) is a small parameter, \( r_1, r_2 \) are defined in Theorem 2 and

\[
P_m(x, y) = \sum_{i=1}^{[r_1/p]} a_i x^{q_i} y^{r_1+i}, \quad Q_n(x, y) = \sum_{j=1}^{[r_2/q]} b_j x^{r_2-jq} y^p.
\]

**Proof.** Let \( l_1 = (r_1 + 1)/2, \ l_2 = (r_2 + 1)/2 \). System (11) is a perturbed system of \((l_1, l_2)\)-quasi-homogeneous Hamiltonian system

\[
\frac{dx}{dt} = y^{r_1}, \quad \frac{dy}{dt} = -x^{r_2}.
\]

It is well known (see for instance [18]) that each simple zero of Abelian integral

\[
I(\rho) = \int_{\Gamma_{\rho}} -P_m(x, y)dy + Q_n(x, y)dx,
\]

produces a periodic orbit of system (11), where \( \Gamma_{\rho} = \{(x, y) : H(x, y) = l_1 x^{2l_2} + l_2 y^{2l_1} = \rho^{2l_1 l_2}\} \) is a closed orbit of system (13). Taking the change \( \rho \), it follows from Lemma 6 and Lemma 7 that

\[
I(\rho) = \sum_{i=1}^{[r_1/p]} a_i \rho^{i(l_1 q - l_2 p) + l_1 l_2} \int_0^T C s^{i q + 2 l_2 - 1} \phi \sin r_1 \in \phi d\phi
\]

\[
+ \sum_{j=1}^{[r_2/q]} b_j \rho^{-j(l_1 q - l_2 p) + r_2 l_1 + l_1} \int_0^T C s^{r_2 - j q} \sin \phi d\phi
\]

\[
= \sum_{i=1}^{[r_1/p]} a_i \rho^{i(m-n)/2 + 2 l_1 l_2} \int_0^T C s^{i q + 2 l_2 - 1} \phi \sin r_1 \in \phi d\phi
\]

\[
+ \sum_{j=1}^{[r_2/q]} b_j \rho^{-j(m-n)/2 + 2 l_1 l_2} \int_0^T C s^{r_2 - j q} \sin \phi d\phi
\]

\[
= \rho^{2 l_1 l_2} \left( \sum_{i=1}^{[(r_1/p) + 1]/2} \mu_{2i-1} \rho^{(2i-1)(m-n)/2} + \sum_{j=1}^{[(r_2/q) + 1]/2} \nu_{2j-1} \rho^{-(2j-1)(m-n)/2} \right),
\]

where

\[
\mu_i = a_i \int_0^T C s^{i q + 2 l_2 - 1} \phi \sin r_1 \in \phi d\phi, \quad \nu_j = b_j \int_0^T C s^{r_2 - j q} \sin \phi d\phi.
\]
If $m \leq n$, then

$$I(\rho) = \rho^{2l_1l_2 + 2[(\lceil r_1/p \rceil + 1)/2] - 1} (m-n)/2 F(\rho^{n-m}),$$

where

$$F(\xi) = \sum_{i=1}^{\lceil (\lceil r_1/p \rceil + 1)/2 \rceil} \mu_{2i-1}^{\lceil (\lceil r_1/p \rceil + 1)/2 \rceil - i} \xi^{\lceil (\lceil r_1/p \rceil + 1)/2 \rceil - 1} + \sum_{j=1}^{\lceil (\lceil r_2/q \rceil + 1)/2 \rceil} \nu_{2j-1}^{\lceil (\lceil r_2/q \rceil + 1)/2 \rceil - 1} \xi^{\lceil (\lceil r_2/q \rceil + 1)/2 \rceil - 1}.$$ 

Since $F(\xi)$ is a polynomial of $\xi$ with degree at most $\lceil (\lceil r_1/p \rceil + 1)/2 \rceil + \lceil (\lceil r_2/q \rceil + 1)/2 \rceil - 1$, $I(\rho)$ has at most $\lceil (\lceil r_1/p \rceil + 1)/2 \rceil + \lceil (\lceil r_2/q \rceil + 1)/2 \rceil - 1$ zeros in $(0, \infty)$. By a suitable choice of $\mu_{2i-1}$’s and $\nu_{2j-1}$’s, we can find a perturbation such that $I(\rho)$ has exactly $\lceil (\lceil r_1/p \rceil + 1)/2 \rceil + \lceil (\lceil r_2/q \rceil + 1)/2 \rceil - 1$ zeros in $(0, \infty)$.

If $m > n$, then result follows by the same arguments as above. □

**Proof of Theorem 2(iii).** Suppose that (ii.1) holds. Consider the $(p, q)$ semi-quasi-homogeneous system (13). Note System (13) is also a $(l_1, l_2)$-quasi-homogeneous Hamiltonian system with a first integral $H(x, y) = l_1 x^{2l_2} + l_2 y^{2l_1}$. The origin is a center of system (13) and all solutions are periodic. This proves (a).

If $b_0 = 0$, then $y = 0$ is an invariant line of system (2) and system (2) has no periodic orbit in this case. This proves (b).

Other assertions in Theorem 2(iii) follow from Theorem 13. □

**Proof of Theorem 2(iv).** Suppose that $L$ is a limit cycle of system (2). Let $\varphi_1(0, y_0)$ be a orbit such that $L$ is a subset of $\omega$ (or $\alpha$) limit set of $\varphi_1(0, y_0)$ and $\varphi_0(0, y_0) = (0, y_0)$. By Poincaré-Bendixson theorem $\omega(\varphi_1(0, y_0)) = L$ (or $\alpha(\varphi_1(0, y_0)) = L$). This implies that there exists $t_1 \in \mathbb{R}$ satisfying $\varphi_1(t_1, 0, y_0) = (0, y_0)$. On the other hand, if $p$ is odd and $q$ is even, then $P_m(-x, y) = P(x, y)$. $Q_n(-x, y) = -Q_n(x, y)$, which yields that the phase portraits of system (2) are symmetric with respect to $y$-axis. Therefore $\varphi_1(0, y_0)$ is a periodic orbit. This shows that $L$ is not a limit cycle (the isolated periodic orbit). Since there is no other singular point in the finite plane, the origin is a center of system (2). □

4. **Proof Theorem 3-5**

In this section we will prove Theorem [3-5].

**Proof of Theorem 3** Without loss of generality suppose $a_0 < 0, b_0 > 0$. Taking the changes

$$\bar{x} = l_0^{1/(r_2+1)} x, \quad \bar{y} = (-a_0)^{1/(r_1+1)} y, \quad \tau = (-a_0)^{1/(r_1+1)} b_0^{1/(r_2+1)} t,$$

system (3) becomes

$$\frac{d\bar{x}}{d\tau} = -\bar{y}^{r_1} + \sum_{i=1}^{\lfloor r_1/p \rfloor} \bar{a}_i \bar{x}^i \bar{y}^{r_1-i}, \quad \frac{d\bar{y}}{d\tau} = \bar{x}^{r_2},$$

where $\bar{a}_i = a_i (-a_0)^{i/(r_1+1) - 1} b_0^{-(q_1 + 1)/(r_2+1)}$. Let $r_1 = 2l_1 - 1$ and $r_2 = 2l_2 - 1$. Note that $-\bar{y}^{r_1} \partial / \partial \bar{x} + \bar{x}^{r_2} \partial / \partial \bar{y}$ is a $(l_1, l_2)$-quasi-homogeneous vector field of weighted degree $2l_1 l_2 - l_2 - l_2 + 1$, and

- (i) If $\bar{a}_1 = \bar{a}_3 = \cdots = \bar{a}_\kappa = 0$, then origin is a center of system (14).
To prove this assertion, we rewrite system (15) in \((\rho, \phi)\) coordinates by \((l_1, l_2)\) polar coordinates introduced in (6):

\[
\frac{d\rho}{d\tau} = \rho^{2l_1l_2-1-l_2+1} \sum_{i=1}^{[r_1/p]} \rho^{(m-n)/2} a_i \cos^{i+2l_2-1} \phi \sin^{r_1-ip+1} \phi,
\]

\[
\frac{d\phi}{d\tau} = \rho^{2l_1l_2-1-l_2} \left( 1 - l_2 \sum_{i=1}^{[r_1/p]} \rho^{(m-n)/2} a_i \cos^{i+2l_2-1} \phi \sin^{r_1-ip+1} \phi \right),
\]

which implies that

\[
(16) \quad \frac{1}{\rho^{1+(m-n)/2}} \frac{d\rho}{d\phi} = \frac{\sum_{i=1}^{[r_1/p]} \rho^{(i-1)(m-n)/2} a_i \cos^{i+2l_2-1} \phi \sin^{r_1-ip} \phi}{1 - l_2 \sum_{i=1}^{[r_1/p]} \rho^{(m-n)/2} a_i \cos^{i+2l_2-1} \phi \sin^{r_1-ip+1} \phi} \leq \frac{A(\rho, \phi)}{G(\rho, \phi)}.
\]

Recall that \(C_s \phi\) and \(S_n \phi\) are \(T\)-periodic functions, where \(T\) is defined in Lemma [4]. For \(\phi \in [0, T]\) and small \(\rho\), we have \(G(\rho, \phi) > 1/2\). This yields that

\[
\left| \frac{A(\rho, \phi)}{G(\rho, \phi)} \right| < M < +\infty,
\]

for \(\phi \in [0, T]\) as \(\rho \to 0\), where \(M\) is a non-negative constant. Integrating both side of \((16)\), one gets

\[
(17) \quad |\rho_2^{-(m-n)/2} - \rho_1^{-(m-n)/2}| = \left| \int_{\phi_1}^{\phi_2} \frac{A(\rho, \phi)}{G(\rho, \phi)} d\phi \right| < M|\phi_2 - \phi_1| < +\infty,
\]

where \(0 \leq \phi_1 < \phi_2 \leq T\) and \(\rho(\phi_i) = \rho_i, \ i = 1, 2\). The above inequality shows that \(\rho(\phi_2) \neq 0\) if \(\rho(\phi_1) \neq 0\). Therefore the origin is either a focus or a center.

Let \(X = (\tilde{P}_m, \tilde{Q}_n)\) be the vector field associated system (15). If \(a_1 = a_3 = \cdots = a_k = 0\), then \(P_m(-x, y) = P(x, y), Q_m(-x, y) = -Q_m(x, y)\), which shows that the phase portraits of system (15) are symmetric with respect \(-x\)-axis. Since the origin is either a focus or a center, the assertion in this part follows.

(ii) If \(m > n\), then \(-\tilde{x}^q \tilde{y}^{r_1-ip} \partial \tilde{x} / \partial \tilde{x}\) is a \((l_1, l_2)\)-quasi-homogeneous vector filed of weighted degree \(2l_1l_2 - l_2 + 1 + i(m-n) > 2l_1l_2 - l_1 - l_2 + 1\) for \(i \geq 1\). Furthermore, if \(i\) is odd, then it follows from Lemma 6 that

\[
I_i = \int \int_{l_1 \tilde{x}^{2l_2} + l_2 \tilde{y}^{2l_1} \leq 1} \text{div}(-\tilde{x}^q \tilde{y}^{r_1-ip}, 0) d\tilde{x} d\tilde{y} = -\int \int_{l_1 \tilde{x}^{2l_2} + l_2 \tilde{y}^{2l_1} = 1} \tilde{x}^q \tilde{y}^{r_1-ip} d\tilde{x} d\tilde{y} = -\int_0^T C_s \tilde{x}^{iq+2l_2-1} \tilde{y}^{r_1-ip} \phi d\phi \neq 0.
\]

The results of this theorem follow from Theorem 9. \(\square\)
Proof of Theorem 4. If \( m = 1 \), then \( r_1 = p/q \). Noting that convention (i) holds, we have \( q = 1 \). On the other hand \( p|q + n - 1 \) and \( q = 1 \) imply that \( p|n \). Therefore \( r_2 = n/p \) is odd integer. System (4) follows from (2).

(i) Suppose \( a_0 = 0, a_1 \neq 0 \). Since \((P_1(x, y), Q_n(x, y)) = 1\), one gets \( b_{n/p} \neq 0 \).

The other assertion of (i) follows from Theorem 8.

(ii) Suppose \( a_0 \neq 0, a_1 \neq 0 \). Assume \( p = 1 \). By means of change \( \bar{x} = -a_1 x - a_0 y, \bar{y} = -a_0 y \), system (4) becomes

\[
\begin{align*}
\frac{d\bar{x}}{dt} &= a_1 \bar{x} + a_0 \sum_{l=0}^{n} b_j \left( \frac{\bar{x} - \bar{y}}{a_1} \right)^{n-j} \left( \frac{\bar{y}}{a_0} \right) \right) \frac{d\bar{y}}{dt} &= a_0 \sum_{l=0}^{n} b_j \left( \frac{\bar{x} - \bar{y}}{a_1} \right)^{n-j} \left( \frac{\bar{y}}{a_0} \right) \right).
\end{align*}
\]

Using Theorem 8 again, the assertion (a.1) and (a.2) follow from (18).

If \( p \geq 3 \), then we get the assertion (b) by Theorem 8.

(iii) The assertion (iii) follows from Theorem 3.

\[\Box\]

Proof of Theorem 5. By the same arguments as in the proof of Theorem 3, the results follow from Remark 10. \[\Box\]

References

[1] D. V. Anosov and V. I. Arnold, Dynamical systems I, Encyclopaedia of Mathematical Sciences 1, Springer-Verlag, Berlin, Heidelberg, New York, 1988.

[2] M. Brunella and M. Miari, Topological equivalence of a plane vector field with its principal part defined through Newton poliedron, J. Differential Equations 85 (1990), 338–366.

[3] L. Cairó, J. Llibre, Phase portraits of planar semi-homogeneous vector fields (I), Nonlineare Analysis 29 (1997), 783–811.

[4] L. Cairó, J. Llibre, Phase portraits of planar semi-homogeneous vector fields (II), Nonlineare Analysis 39 (2000), 338–366.

[5] A. Cima, A. Gasull and F. Mañàs, Cyclicity of a family of vector fields, J. Math. Anal. Appl. 196 (1995), 420–448.

[6] A. Cima, A. Gasull and F. Mañàs, Limit cycles for vector fields with homogeneous components, Applicationes Mathematicae 24(1997), 281–287.

[7] A. Cima and J. Llibre, Algebraic and topological classification of homogeneous cubic vector fields in the plane, J. Math. Anal. Appl. 147 (1990), 420–448.

[8] B. Coll, A. Gasull and P. Prohens, Differential equations defined by the sum of two quasi-homogeneous vector fields, Can. J. Math. 13(1996), 63–91.

[9] J. Chavarriga and I. A. García, Lie symmetries of quasihomogeneous polynomial planar vector fields and certain perturbations, Acta Math. Sin. (Engl.Ser.) 21 (2005), 185–192.

[10] F. Dumortier, J. Llibre, J. C. Artés, Qualitative theory of planar differential systems, Springer, 2006.

[11] L. Gavrilov, J. Giné and Maite Grau, On the cyclicity of weight-homogeneous centers centers, J. Differential Equations 246(2009), 3126-3135.

[12] A.M. Liapunov, Stability of motion, in Mathematics in Science and Engineering, Vol. 30, Academic Press, 1966.
[13] Weigu Li, J. Llibre, J. Yang and Z. Zhang, *Limit cycles bifurcating from the period annulus of quasi-homogeneous centers*, J. Dynam. Differential Equations 21(2009), 133-152.

[14] J. Llibre, Jesus S. Perez del Rio and J.A. Rodriguez, *Structural stability of planar homogeneous polynomial vector fields: applications to critical points and to infinity*, J. Differential Equations 125 (1996), 490–520.

[15] J. Llibre, Jesus S. Perez del Rio and J.A. Rodriguez, *Structural stability of planar semi-homogeneous polynomial vector fields: applications to critical points and to infinity*, Discrete and Continuous Dynamical Systems, 6 (2000), 809–828.

[16] J. Llibre and X. Zhang, *Polynomial first integrals for quasi-homogeneous polynomial differential systems*, Nonlinearity, 15 (2002), 1269–1280.

[17] R.D.S.Oliveira and Yulin Zhao, *On the structural stability of planar quasi-homogeneous polynomial vector fields*, Preprint.

[18] R. Roussarie, *Bifurcation of planar vector fields and Hilbert’s sixteenth problem*, Progress in Mathematics 164, Birkhäuser Verlag, Basel, 1998.

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