An Iterative Algorithm for the Estimation of Thermal Boundary Conditions Varying in Both Time and Space

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Abstract: The presented survey of the up-to-date state of knowledge indicates that despite the great number of works devoted to the issue in question, there is no simple method that allows the use of commercial programs for the identification of the transient thermal state in elements with a simple or complex shape. This paper presents an inverse method developed to estimate the convective heat transfer coefficient varying both in time and space on a vertical plate during its cooling. Despite the smaller number of measurement points and larger disturbance of measured temperatures compared to the data presented in the available literature, comparable results are obtained. The developed iterative algorithm is also applied to estimate the time- and space-dependent heat flux and the convective heat transfer coefficient in the steam boiler membrane waterwall. The analysed component has the form of the non-simply connected and complex shape domain $\Omega$. Temperature-dependent thermophysical properties are used. Calculations are performed for the unknown heat flux or heat transfer coefficient distribution on the domain boundary based on measured temperature transients disturbed with a random error of 0.5 °C. To reduce oscillations, the number of future time steps of $N_F = 20$ is selected. The number of iterations in each time step ranges between 1 and 8. The estimated boundary conditions are close to the exact values. In this work, the ANSYS software using the FEM is applied.

Keywords: convective heat transfer coefficient estimation; inverse heat conduction problem; nonlinear estimation; Levenberg-Marquardt algorithm; steam boilers

1. Introduction

Modelling heat and fluid flows as well as thermal-, flow-, and strength-related processes plays an important role in the design and operation in engineering. The influence of the real ground temperature on a building’s heat loss is analysed in [1]. The numerical model of heating and cooling unit with a reversible heat pump based on CO$_2$ is presented in [2]. Its experimental validation is also reported. Paper [3] presents a novel crossflow air-to-water heat exchange, its experimental testing, and modelling. A thermal and strength analysis of the power boiler superheater is shown in [4]. The maximum operating temperature is calculated by the proposed creep equation. The transient-state distribution of thermal stresses in a steam gate valve and the proposition of optimization algorithm for heating and cooling operations is presented in [5]. The significant problem of modelling the temperature and stress distributions in newly designed and used components causes the difficulty in defining some of the boundary conditions. Modelling requires the specification of dimensions, thermophysical properties, source terms, boundary conditions or initial conditions. The more precisely the input data are determined, the better the modelling results are. In practice, it is often difficult to define the boundary conditions. Placing thermocouples on the inner surface of CFB furnace, superheater chambers, pipelines, as well as gates and valves can destroy them [6]. A direct temperature measurement on the outer surface of an atmospheric re-entry capsule is difficult due to its high value [7].

A boundary condition that is difficult to define can be considered as unknown, and additional information can be added to the analysis. An analysis defined in this way is...
called an inverse boundary problem. For components with simple shapes, one-dimensional temperature distribution and constant material properties’ exact methods can be used [8]. A certain inconvenience of the proposed solutions is the need to take a larger number of terms in the series describing the temperature distribution. If the number of terms is not high enough, the temperature distribution calculated at the initial time instant will differ from the real one. For multidimensional problems, finding a solution is laborious [9]; therefore, numerical methods are used. Nonlinear one-dimensional inverse problems may be solved by means of the finite difference method (FDM) [10]. The proposed method for a two-dimensional problem [11] has a global character, which means that calculations may be started only after measurements are performed, when the temperature histories from the initial to the final time instant have been recorded. For this reason, the methods cannot be used for online identification of the temperature distribution in the element. An optimization method intended for solving multidimensional inverse problems was proposed in [12]. In [13], an online inverse space marching method based on the control volume method (CVM) in cylindrical coordinates was formulated. The proposed method was applied to identify the transient-state two-dimensional temperature estimation in the steam header cross-section and to reconstruct the unknown boundary condition on its inner surface. A numerical algorithm for determination of the heat transfer coefficient distribution on a vertical plate under mixed-convection conditions was developed in [14] using a two-dimensional inverse method. Applications of a modified Levenberg-Marquardt method for identification of the conductivity and heat capacity of solids is presented in [15,16].

The presented survey of the up-to-date state of knowledge indicates that despite the great number of works devoted to the issue in question, there is no simple method that allows the use of commercial programs for the identification of the transient thermal state in elements with a simple or complex shape. The methods published so far require writing the source code with the use of FDM or CVM.

This paper presents a method developed to estimate the convective heat transfer coefficient varying both in space and time. The proposed method is used to estimate the heat transfer coefficient distribution of a vertical plate during cooling under the mixed-convection regime. The estimated values are compared to those calculated in [6]. Despite the smaller number of measurement points and larger disturbance of the measured temperatures, the obtained results are comparable to [14]. Additionally, the proposed method is easier to apply as it allows the use of commercial programs. The developed iterative algorithm is also applied to estimate the heat flux and the convective heat transfer coefficient varying over time in the steam boiler membrane waterwall. The analysed component has the form of the non-simply connected domain \( \Omega \).

2. Formulation of the Method

The transient heat conduction problem in the non-simply connected domain \( \Omega \) shown in Figure 1 can be described using the following equation [17]:

\[
c(T) \rho(T) \frac{\partial T}{\partial t} = -\nabla \cdot \mathbf{q}
\]

(1)

where \( \mathbf{q} \) is the heat flux vector defined by Fourier’s law:

\[
\mathbf{q} = -D \nabla T
\]

(2)

and

\[
D = \begin{bmatrix}
k_x(T) & 0 & 0 \\
0 & k_y(T) & 0 \\
0 & 0 & k_z(T)
\end{bmatrix}
\]

(3)

If the material is isotropic, then \( k_x(T) = k_y(T) = k_z(T) = k(T) \). All material properties (\( c \)—specific heat, \( k \)—thermal conductivity, and \( \rho \)—density) are usually known functions of temperature. Fourier’s law applies to all engineering applications for which the infinite
speed of a heat propagation can be assumed. Transient heat conduction problems are initial–boundary problems for which appropriate initial and boundary conditions should be defined [17]. The initial condition specifies the temperature value of the non-simply connected domain \( \Omega \) at its first moment \( t_0 = 0 \) s:

\[
T(x, y, z, 0)|_{t_0=0} = T_0(x, y, z)
\]  

(4)

One of the first-, second-, or third-kind boundary conditions should be defined on boundary \( \Gamma = \Gamma_{in} \cup \Gamma_{out} \):

\[
T|_{\Gamma} = T_b
\]  

(5)

\[
(\mathbf{D} \nabla T \cdot \mathbf{n})|_{\Gamma} = q_B
\]  

(6)

\[
(\mathbf{D} \nabla T \cdot \mathbf{n})|_{\Gamma} = h(T_m - T|_{\Gamma})
\]  

(7)

where \( \mathbf{n} \)—unit normal vector to boundary, \( T_b \)—temperature on the domain boundary, \( q_B \)—heat flux on the domain boundary, \( h \)—heat transfer coefficient on the domain boundary \( \Gamma \), and \( T_m \)—temperature of the medium. A phenomenon characterized by Equations (1)–(7) in which all dimensions, properties, the initial condition, and the boundary conditions over the entire boundary \( \Gamma = \Gamma_{in} \cup \Gamma_{out} \) are specified is known as a direct transient heat conduction problem. If part of boundary conditions (5)–(7) is unknown, the problem is ill-posed, and additional internal temperature measurements are required

\[
f_i(t) = T(r_i) \quad i = 1, \ldots, N_T
\]  

(8)

Figure 1. Non-simply connected domain \( \Omega \) bounded by boundary \( \Gamma = \Gamma_{in} \cup \Gamma_{out} \).

The unknown first-, second-, or third-kind boundary condition is divided into \( N_u \) intervals. It can be approximated as a staircase or a polynomial function on the domain edge (Figure 2).

Figure 2. Staircase or polynomial approximation of the unknown boundary condition on the domain edge.
The aim is to choose time-varying $u(t_1(t), u_2(t), \ldots, u_{N_u}(t))$ such that the calculated temperatures are close to the temperature values obtained from the experimental measurements. This can be expressed as:

$$T_i(u, r_i, t) - f_i(t) \approx 0, \quad i = 1, \ldots, N_T$$

(9)

The number of measurement points $N_T$ cannot be smaller than the number of unknown boundary condition intervals $N_u$. Ill-posed problems are tough to solve. They are sensitive to errors in input data. Many algorithms can be adopted to stabilise the solution. Future time steps are applied in this paper. To stabilise the solution, vector $u_L$ is held constant in $N_F$ future time steps [18], as can be seen in Figure 3.

$$u_L = u_{L+1} = \ldots = u_{L+N_F-1}$$

(10)

where $u_L = u(u_1(t_L), u_2(t_L), \ldots, u_{N_u}(t_L))$.

![Diagram illustrating the calculation of $u_L$ using future time steps.](image)

Figure 3. Diagram illustrating the calculation of $u_L$ using future time steps.

The least squares method is applied to estimate parameters $u_L$. The sum of squares

$$S = \sum_{i=1}^{N_T} \sum_{j=L}^{L+N_F-1} \left[ f(t_j) - T(u_L, r_i, t_j) \right]^2$$

(11)

can be minimized using a general unconstrained method. The Levenberg-Marquardt method [15] is applied herein to determine parameters $u_L$. It is assumed that the sought vector changes in the range of:

$$u_L^l \leq u_L \leq u_L^u$$

(12)

where superscripts $l$ and $u$ denote the lower and the upper bound of the sought parameters, respectively. The initial–boundary problem given by Equations (1)–(7) is solved in each iteration step by the finite element method (FEM) calculating the unknown boundary conditions from $u_L$.

The iterative algorithm for the inverse problem solution using the LM method can be summarized as follows:

1. Solve the initial–boundary problem formulated by Equations (1)–(7) in the interval $t \in [0, \Delta t N_F]$ for the initial value of sought vector $u_L(L = 0)$ from the range defined by (12);
2. Determine objective function $S$ using (11);
3. Apply the LM method to estimate vector \( \mathbf{u}_L(L = 0) \), which minimizes objective function \( S \);  
4. Save the obtained vector \( \mathbf{u}_L(L = 0) \) and the temperature distribution in time \( t = \Delta t N_F \).  
5. Solve the initial–boundary problem formulated by Equations (1)–(7) in interval \( t \in \langle \Delta t N_F, 2\Delta t N_F \rangle \) taking the initial value of vector \( \mathbf{u}_L(L = N_F) = \mathbf{u}_L(L = 0) \) and the obtained temperature distribution in time \( t = \Delta t N_F \) for the initial condition (4);  
6. Apply the LM method to estimate vector \( \mathbf{u}_L(L = N_F) \), which minimizes objective function \( S \);  
7. Save vector \( \mathbf{u}_L(L = N_F) \) and the temperature distribution in time \( t = 2\Delta t N_F \);  
8. Continue the algorithm to the time \( (t_{end}) \).

The proposed algorithm is used to identify the heat transfer coefficient distribution of a heated vertical plate during cooling under the mixed-convection regime. The applicability of the proposed algorithm will also be demonstrated by solving an inverse heat conduction problem in the non-simply connected domain \( \Omega \). It will be used for a transient-state analysis of membrane waterwalls in steam boilers.

3. Mixed Convection on a Vertical Plate

A plate with initial temperature \( T_0 = 200 \, ^\circ C \), height \( L = 20 \) cm, and width \( b = 1 \) cm is presented in Figure 4 [14]. It was insulated on three sides and vertically exposed to mixed convection from the left surface. The material properties of the stainless-steel plate in the adopted initial temperature were as follows: \( \rho = 8238 \, [kg/m^3] \), \( k = 15.2 \, [W/mK] \), and \( c_p = 504 \, [J/kg\cdot K] \).

![Figure 4. A plate exposed to mixed convection.](image)

Combined free and forced convection affect a vertical hot plate due to an assisting flow with temperature \( T_\infty = 25 \, ^\circ C \), while the Richardson number \( Ri = Gr_L/Re_L^2 \) varies from 0.1 to 1. All fluid properties were calculated at the mean value of film temperature. The mixed-convection heat transfer was then evaluated along the plate \( Nu = (Nu_F^3 + Nu_N^3)^{1/3} \), where \( Nu_F \) is the local Nusselt number of forced convection [19], and \( Nu_N \) is the local Nusselt number of natural convection [20]. Finally, the convective heat transfer varying both over time and along the plate length on the left boundary \( (h_{mc}(x, t)) \) was calculated. The heat transfer distribution in space was approximated using the following function

\[
h_{mc}(x, t) = a(t) + b(t)/x^{0.5}
\]

where parameters \( a \) and \( b \) were found by the least squares method in every minute of cooling. They are presented in Table 1.
Table 1. Coefficients of the function approximating the distribution of the convective heat transfer in time and space in Equation (13).

| \(t\) [s] | \(a\) [W/m\(^2\)K] | \(b\) [W/m\(^3\2\)K] |
|----------|----------------|----------------|
| 60       | 1.2198462     | 18.67917202   |
| 120      | −0.162362     | 25.376006     |
| 180      | −1.194652     | 31.048731     |
| 240      | −2.005685     | 35.961097     |
| 300      | −2.671        | 40.32         |
| 360      | −3.238473     | 44.280206     |
| 420      | −3.734243     | 47.914765     |
| 480      | −4.181083     | 51.307252     |
| 540      | −4.581333     | 54.476993     |
| 600      | −4.955758     | 57.48656938   |

The temperature distribution during the cooling process was calculated using the FEM implemented in the ANSYS program [21]. ANSYS Mechanical APDL was used in the batch mode. The plate was discretized into 2000 linear plane finite elements. Ten elements were used along its thickness and 200 along its height. The convective boundary condition was assumed on the respective surfaces of 200 elements. The temperature histories obtained from the FEM were used as “exact measured data”. When the calculations were carried out with the time step of \(\Delta t = 60\) s, the temperature distribution was very close to the one presented in [14], which was measured and generated by the MATLAB code. Figure 5 shows temperature transients \(f_1(t), f_2(t)\) and \(f_3(t)\) in points \(P_1\), \(P_2\), and \(P_3\), respectively, whose location is given in mm based on the coordinate system in Figure 4: (3, 1), (62, 1), and (191, 1). A reduction in the time step to \(\Delta t = 10\) s involved a slight difference in the temperature transients.

![Figure 5. Comparison between temperature transients from the MATLAB code and from the FEM with the time step of \(\Delta t = 60\) s and \(\Delta t = 10\) s.](image)

The unknown time- and space-dependent heat transfer coefficient of cooling in the mixed-convection regime was calculated by the proposed algorithm based on the three measured temperature transients \(f_1(t), f_2(t),\) and \(f_3(t)\) described in Figure 4.
It was assumed that the unknown parameters $a$ and $b$ from Equation (13) were functions of time. There were two unknown transients and three given measured temperature transients. Vector $\mathbf{u}_L$ consisted of two components: $a(t_L)$ and $b(t_L)$.

The temperature histories taken from the FEM with time step $\Delta t = 10$ s were considered as “exact measured data”. To obtain numerical tests similar to real conditions, “noisy measured data” were formulated by adding normal random errors of $\pm 0.5 \, ^\circ \text{C}$ of the zero mean to “exact measured data” ($\sigma_t = 1/6$). These temperature transients in nodes $P_1$, $P_2$, and $P_3$ for the inverse solution are shown in Figure 6. They are compared to the data presented in the available literature [14].

Figure 6. “Noisy measured data” generated by the FEM for the inverse solution and temperature transients form the MATLAB code.

The presented method was tested using the produced “noisy measured data” and time step $\Delta t = 10$ s. The values of $<-10 \, \text{W/m}^2\text{K}, 0 \, \text{W/m}^{3/2}\text{K}>$ and $<10 \, \text{W/m}^2\text{K}, 100 \, \text{W/m}^{3/2}\text{K}>$, respectively, were chosen in vector $\mathbf{u}_L$ for the lower and the upper bound of all sought parameters. The starting vector was assumed as $\mathbf{u}_L(L = 0) = <1 \, \text{W/m}^2\text{K}, 20 \, \text{W/m}^{3/2}\text{K}>$. Calculations without future time steps could produce instability. To reduce oscillations, the number of future time steps was selected as $N_F = 6$. The number of iterations in each time step ranged between 3 and 30. If the midpoints of the constant parameter intervals were connected with straight lines, the estimated values of $a$ and $b$ were close to the exact ones as can be seen in Figure 7.

Finally, the unknown time- and space-dependent heat transfer coefficient distribution in the mixed-convection regime was calculated. The comparison between estimated and exact values is presented in Figure 8. The HTC transients in Figure 8a are shown for the following $x$ coordinates: 1, 6, 16, 59, and 195 mm. They are also compared to the heat transfer coefficients calculated by the MATLAB code in paper [14]. It needs to be noted that in [14] twenty measurement points were used, and the investigation considered errors of $\pm 0.1 \, ^\circ \text{C}$. Despite the smaller number of measurement points and larger disturbance of measured temperatures than in [14], the obtained results were comparable. Additionally, the proposed method was easier to apply, as it allowed the use of commercial programs.
Figure 6. “Noisy measured data” generated by the FEM for the inverse solution and temperature transients form the MATLAB code.

Figure 7. Comparison between exact parameters and parameters calculated by the inverse method based on “noisy measured data”.

Finally, the unknown time- and space-dependent heat transfer coefficient distribution in the mixed-convection regime was calculated. The comparison between estimated and exact values is presented in Figure 8. The HTC transients in Figure 8a are shown for the following \( x \) coordinates: 1, 6, 16, 59, and 195 mm. They are also compared to the heat transfer coefficients calculated by the MATLAB code in paper [14]. It needs to be noted that in [14] twenty measurement points were used, and the investigation considered errors of \( \pm 0.1 \) °C. Despite the smaller number of measurement points and larger disturbance of measured temperatures than in [14], the obtained results were comparable. Additionally, the proposed method was easier to apply, as it allowed the use of commercial programs.

The introduced random errors had a slight effect on the identified boundary condition. Future steps effectively eliminated random measurement errors. During each iteration, the initial–boundary problem formulated by Equations (1)–(7) was solved by the FEM using the ANSYS software [21]. Due to the parametric formulation of the method, its use was simplified. The presented algorithm can be applied to transient nonlinear problems if they only have a direct solution.

4. Transient Analysis of Membrane Waterwalls in Steam Boilers

The furnace wall tubes are usually welded together using steel bars to make membrane wall panels, which are exposed to the furnace on one side and insulated on the other, as presented in Figure 9. A tubular-type instrument (flux tube) was installed to enable accurate measurement of the absorbed heat flux \( q_m \) and heat transfer coefficient \( h \) inside the membrane wall panels.
Due to symmetry, only the representative waterwall section presented in Figure 11 needs to be analysed. Temperature transients inside the membrane wall in steady-state conditions [22]. The proposed method was used to identify the absorbed heat flux and the heat transfer coefficient during transient operation.

Temperature-dependent material properties were assumed (Figure 10). The following functions of the absorbed heat flux and of the heat transfer coefficient over time were assumed:

\[
\begin{align*}
q_m(t) &= 14.03t + 39,038.8 \quad [\text{W/m}^2] \\
h(t) &= -1.12t + 14,970.9 \quad [\text{W/m}^2\text{K}] \\
q_m(t) &= -1.025t + 53,068.2 \quad [\text{W/m}^2] \\
h(t) &= 8.66t + 13,849.5 \quad [\text{W/m}^2\text{K}] \\
\end{align*}
\]

for \(0 < t < 1000\) s

for \(1000 < t < 2000\) s

Figure 9. Cross-section of the membrane wall with a flux tube in the steam boiler combustion chamber: 1—waterwall tube, 2—flux tube, \(f_1(t)–f_5(t)\)—location of thermocouples to measure temperature.

The membrane wall and the flux tube were made of 20 G steel. The furnace wall tubes are usually welded together using steel bars to make nonlinear problems if they only have a direct solution.

The membrane wall and the flux tube were made of 20 G steel. Temperature-dependent material properties were assumed (Figure 10). The proposed method was used to identify the absorbed heat flux and the heat transfer coefficient during transient operation.

Steam temperature totals \(T_m = 316.52\) [°C]. The local heat flux on the membrane wall and on the flux tube was evaluated numerically using the FEM-based ANSYS program [21]. Due to symmetry, only the representative waterwall section presented in Figure 11 needs to be analysed.
to be analysed. Temperature transients $f_1, f_2, f_3, f_4,$ and $f_5$ in the locations illustrated in Figure 9 were calculated using the FEM and disturbed with normal random errors ±0.5 °C of the zero mean ($\sigma_f = 1/6$). The generated “noisy measured data” for the inverse solution are shown in Figure 12.

Figure 11. Division of the non-simply connected domain into finite elements: 1—flux tube, 2—waterwall tube, $f_1(t)$—$f_5(t)$—location of temperature sensors.

Figure 12. “Noisy measured data” generated by the FEM for the inverse solution.

There were two unknown transients $q_m(t), h(t)$ and five given measured temperature transients. Vector $\mathbf{u}_L$ consisted of two components: $u_1(t_L)$ and $u_2(t_L)$.

The proposed method was tested with the time step of $\Delta t = 10$ s. The values of 0 and 100,000 <W/m² K, W/m²>, respectively, were chosen in vector $\mathbf{u}_L$ for the lower and the upper bounds of the sought parameters. The starting vector was assumed as $\mathbf{u}_L(L = 0) = <$15,000 W/m² K, 40,000 W/m²$>

To reduce oscillations, the number of future time steps of $N_F = 20$ was selected. The number of iterations in each time step ranged between 1 and 8. The estimated boundary conditions were close to the exact values, as can be seen in Figure 13.

Figure 13. Comparison between exact and identified boundary conditions based on “noisy measured data”.

The membrane wall and the flux tube were made of 20 G steel. Temperature-dependent material properties were assumed (Figure 10). The following functions of the absorbed heat flux and of the heat transfer coefficient over time were assumed:

$$
q(t) = 14.03t + 39038.8 \text{ W/m}^2 \text{ for } 0 \leq t \leq 1000 \text{ s} \\
h(t) = 1.12t - 1.025t^2 + 53068.2 \text{ W/m}^2 \text{ K} \text{ for } 1000 \leq t \leq 2000 \text{ s}
$$

Steam temperature totals $T_m = 316.52$ °C. The local heat flux on the membrane wall and on the flux tube was evaluated numerically using the FEM-based ANSYS program [21]. Due to symmetry, only the representative waterwall section presented in Figure 11 needs to be analysed. Temperature transients $f_1, f_2, f_3, f_4,$ and $f_5$ in the locations illustrated in Figure 9 were calculated using the FEM and disturbed with normal random errors ±0.5 °C of the zero mean ($\sigma_f = 1/6$). The generated “noisy measured data” for the inverse solution are shown in Figure 12.

Figure 10. Material properties of 20 G steel.
There were two unknown transients \( q_m(t), h(t) \) and five given measured temperature transients. Vector \( u_L \) consisted of two components: \( u_1(t) \) and \( u_2(t) \).

The proposed method was tested with the time step of \( \Delta t = 10 \text{s} \). The values of 0 and 100,000 \( \text{W/m}^2 \text{K} \), \( \text{W/m}^2 > \), respectively, were chosen in vector \( u_L \) for the lower and the upper bounds of the sought parameters. The starting vector was assumed as \( u_L(L) = 0) = <15,000 \text{W/m}^2 \text{K}, 40,000 \text{W/m}^2> \).

To reduce oscillations, the number of future time steps of \( N_F = 20 \) was selected. The number of iterations in each time step ranged between 1 and 8. The estimated boundary conditions were close to the exact values, as can be seen in Figure 13.

Figure 13. Comparison between exact and identified boundary conditions based on “noisy measured data”.

5. Conclusions

The presented survey of the up-to-date state of knowledge indicated that despite the great number of works devoted to the issue in question, there was no simple method, which allowed the use of commercial programs for the identification of the transient thermal state in elements with a simple or complex shape. The proposed method did not require writing the source code with the use of FDM, CVM, or FEM. Its advantage was the possibility of using commercial programs. In this work, the ANSYS software using the FEM was applied. The two presented examples demonstrated good stability and accuracy of the developed method. The proposed inverse method made it possible to estimate the convective heat transfer coefficient in the heated vertical plate varying both in time and space during cooling under the mixed-convection regime. Despite the smaller number of measurement points and greater disturbance of measured temperatures than in [14], results comparable to [14] were obtained. The developed iterative algorithm was also applied to estimate the time-dependent heat flux and the convective heat transfer coefficient in the steam boiler membrane waterwall. The analysed component had the form of the non-simply connected domain \( \Omega \), and temperature-dependent thermophysical properties were assumed. Calculations were performed for the unknown heat flux or the heat transfer coefficient distribution on the domain boundary based on measured temperature transients disturbed with the random error of 0.5 \( ^\circ \text{C} \). To reduce oscillations, the number of future time steps of \( N_F = 20 \) was selected. The number of iterations in each time step ranged between 1 and 8. The estimated boundary conditions were close to the exact values. The method can be used in practice in an offline mode.

Author Contributions: Conceptualization, P.D.; methodology, P.D.; software, M.K.; validation, M.K.; formal Analysis, P.D.; investigation, P.D.; writing—original draft preparation, P.D.; writing—review & editing, P.D.; visualization, M.K.; supervision, P.D. All authors have read and agreed to the published version of the manuscript.

Funding: This research received no external funding.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.
Conflicts of Interest: The authors declare that there is no conflict of interest regarding the publication of this paper.

Nomenclature

\begin{align*}
c & \text{ specific heat [J/kgK]} \\
k & \text{ thermal conductivity [W/mK]} \\
D & \text{ thermal conductivity matrix [W/mK]} \\
f & \text{ measured temperature history [\degree C]} \\
h & \text{ heat transfer coefficient [W/m\degree K]} \\
N_F & \text{ number of future time steps [-]} \\
N_T & \text{ number of temperature measurement points [-]} \\
N_u & \text{ number of unknown boundary condition intervals [-]} \\
t & \text{ time [s]} \\
T & \text{ temperature [\degree C]} \\
u & \text{ vector of unknown boundary conditions, individual components [W/m\degree or W/m\degree K]} \\
qu & \text{ heat flux [W/m\degree]} \\
qu_m & \text{ absorbed heat flux [W/m\degree]} \\
q & \text{ heat flux vector [W/m\degree]} \\
\rho & \text{ density [kg/m\degree^3]} \\
\sigma_f & \text{ standard deviation}
\end{align*}

References

1. Pokorska-Silva, I.; Kadela, M.; Orlik-Kożdoń, B.; Fedorowicz, L. Calculation of Building Heat Losses through Slab-on-Ground Structures Based on Soil Temperature Measured In Situ. *Energies* 2022, 15, 114. [CrossRef]
2. Artuso, P.; Tosato, G.; Rossetti, A.; Marinetti, S.; Hafner, A.; Banasiak, K.; Minetto, S. Dynamic Modelling and Validation of an Air-to-Water Reversible R744 Heat Pump for High Energy Demand Buildings. *Energies* 2021, 14, 8238. [CrossRef]
3. Kailkhura, G.; Mandel, R.K.; Shooshhtari, A.; Ohadi, M. Numerical and Experimental Study of a Novel Additively Manufactured Metal-Polymer Composite Heat-Exchanger for Liquid Cooling Electronics. *Energies* 2022, 15, 598. [CrossRef]
4. Duda, P.; Felkowski, L.; Duda, A. An Analysis of Creep Phenomena in the Power Boiler Superheaters. *Metals* 2018, 8, 892. [CrossRef]
5. Duda, P.; Dwornicka, R. Optimization of heating and cooling operations of steam gate valve. *J. Struct. Multidiscip. Optim.* 2010, 40, 529–535. [CrossRef]
6. Blaszczyk, A.; Nowak, W. Bed-to-wall heat transfer coefficient in a supercritical CFB boiler at different bed particle sizes. *Int. J. Heat Mass Transf.* 2014, 79, 736–749. [CrossRef]
7. Duda, P.; Nakamura, T. Identification of the transient temperature and stress distribution in an atmospheric reentry capsule assuming temperature-dependent material properties. *Aerosp. Sci. Technol.* 2017, 67, 265–272. [CrossRef]
8. Burggraf, O.R. An exact solution of the inverse problem in heat conduction theory and applications. *ASME J. Heat Transf.* 1964, 86, 373–382. [CrossRef]
9. Alnajem, N.M.; Özisik, M.N. On the solution of three-dimensional inverse heat conduction in finite media. *Int. J. Heat Mass Transf.* 1985, 28, 2121–2128. [CrossRef]
10. Raynaud, M.; Bransier, J. A new finite difference method for the nonlinear inverse heat conduction problem. *Numer. Heat Transfer* 1986, 9, 30–37. [CrossRef]
11. Busby, H.R.; Trujillo, D.M. Numerical solution to a two-dimensional inverse heat conduction problem. *Int. J. Numer. Methods Eng.* 1985, 21, 349–359. [CrossRef]
12. Alianov, O.M.; Nenarokomov, A.V. Three-dimensional boundary inverse heat conduction problem for regular coordinate systems. *Inverse Probl. Eng.* 1999, 7, 335–362. [CrossRef]
13. Duda, P. Numerical and experimental verification of two methods for solving an inverse heat conduction problem. *Int. J. Heat Mass Transf.* 2015, 84, 1101–1112. [CrossRef]
14. Razzaghi, H.; Kowsary, F.; Ashjaaee, M. Derivation and application of the adjoint method for estimation of both spatially and temporally varying convective heat transfer coefficient. *Appl. Therm. Eng.* 2019, 154, 63–75. [CrossRef]
15. Yang, K.; Jiang, G.H.; Peng, H.F.; Gao, X.W. A new modified Levenberg-Marquardt algorithm for identifying the temperature-dependent conductivity of solids based on the radial integration boundary element method. *Int. J. Heat Mass Transf.* 2019, 144, 118615. [CrossRef]
16. Sassine, E.; Cherifb, Y.; Antczak, E. Parametric identification of thermophysical properties in masonry walls of buildings. *J. Build. Eng.* 2019, 25, 100801. [CrossRef]
17. Özisik, M.N. *Heat Conduction*; John Wiley & Sons: New York, NY, USA, 1980.
18. Beck, J.V. Nonlinear estimation applied to the nonlinear inverse heat conduction problem. *Int. J. Heat Mass Transf.* **1970**, *13*, 703–716. [CrossRef]

19. Bergman, F.P.; Incropera, A.S.; Lavine, D.P. *Introduction to Heat Transfer*; John Wiley & Sons: Hoboken, NJ, USA, 2011.

20. Ostrach, S. *An Analysis of Laminar Free-Convection Flow and Heat Transfer about a Flat Plate Parallel to the Direction of the Generating Body Force*; (No. NACA-TN-2635); National Aeronautics and Space Administration Cleveland Oh Lewis Research Center: Cleveland, OH, USA, 1952.

21. ANSYS User’s Manual, Revision 5.6. Available online: [http://research.me.udel.edu/~jlwang/teaching/MEx81/ansys56manual.pdf](http://research.me.udel.edu/~jlwang/teaching/MEx81/ansys56manual.pdf) (accessed on 5 April 2022).

22. Taler, J.; Duda, P.; Węgłowski, B.; Zima, W.; Grądziel, S.; Sobota, T.; Taler, D. Identification of local heat flux to membrane waterwalls in steam boilers. *Fuel* **2009**, *88*, 305–311. [CrossRef]