3D Hierarchical Rutile TiO₂ and Metal-free Organic Sensitizer Producing Dye-sensitized Solar Cells 8.6% Conversion Efficiency

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Three-dimensional (3D) hierarchical nanoscale architectures comprised of building blocks, with specifically engineered morphologies, are expected to play important roles in the fabrication of 'next generation' microelectronic and optoelectronic devices due to their high surface-to-volume ratio as well as opto-electronic properties. Herein, a series of well-defined 3D hierarchical rutile TiO₂ architectures (HRT) were successfully prepared using a facile hydrothermal method without any surfactant or template, simply by changing the concentration of hydrochloric acid used in the synthesis. The production of these materials provides, to the best of our knowledge, the first identified example of a ledgewise growth mechanism in a rutile TiO₂ structure. Also for the first time, a Dye-sensitized Solar Cell (DSC) combining a HRT is reported in conjunction with a high-extinction-coefficient metal-free organic sensitizer (D149), achieving a conversion efficiency of 5.5%, which is superior to ones employing P25 (4.5%), comparable to state-of-the-art commercial transparent titania anatase paste (5.8%). Further to this, an overall conversion efficiency 8.6% was achieved when HRT was used as the light scattering layer, a considerable improvement over the commercial transparent/reflector titania anatase paste (7.6%), a significantly smaller gap in performance than has been seen previously.

The design and growth of inorganic nanostructures with well-controlled sizes and morphologies have been the focus of intensive research in recent years due to their wide ranging applications. TiO₂ is a functional material used in many ways, such as in photocatalysis, photoelectrochemical solar energy conversion¹, and Li-ion batteries². The properties and applications of TiO₂ are generally correlated to crystal phase, dimensions, morphology, and surface properties. As such, rationalized designed and controlled nanostructured TiO₂ has attracted great interest³,⁴. Various TiO₂ nanostructures in the forms of spheres, tubes, wires, rods, sheets, belts, flowers, and trees have been reported⁵–¹². Three-dimensional (3D) hierarchical structures are of particular interest as many of these architectures have distinctive physicochemical properties in comparison with simple nanocrystals¹³,¹⁴. When employed as the photoanode in a Dye-sensitized Solar Cell (DSC), hierarchical architectures can offer larger surface areas for dye adsorption while simultaneously enhancing light scattering for efficient photon harvesting, and thereby, improving power conversion efficiency¹⁵–¹⁸. Chemical vapor deposition, applications of templates, catalysts and electrochemical methods are known to be used for the direct growth of inorganic material with controlled morphologies and architectures¹⁹–₂⁴, however, in most cases, the mechanisms involved in the formation of these hierarchical structures are complicated and not well-understood.

Crystalline TiO₂ mainly exists in three polymorphs: anatase (tetragonal), rutile (tetragonal) and brookite (orthorhombic). The majority of research has been carried out on anatase, while rutile is often ignored in DSCs, possibly due to the expectation that rutile would result in a lower open-circuit voltage (Voc) than anatase because rutile has a more positive conduction band edge potential.₂⁵ Slow electron transport is also often seen as a detracting feature of rutile TiO₂, this, however, may lead to higher electron densities within the conduction band, subsequently increasing the quasi Fermi-level, and mitigating the aforementioned concern about conduc-
tion band edge potential. Indeed, studies show that rutile TiO₂ can achieve a similar²⁶ or even higher $V_{oc}$ through doping²⁷. Furthermore, compared to anatase, the rutile polymorph exhibits some superior physical properties, such as enhanced light-scattering properties on account of its higher refractive index, which is beneficial from the perspective of effective light harvesting²⁸. Rutile TiO₂ is also chemically more stable and potentially cheaper to produce than anatase²⁹. These findings suggest that additional research on rutile TiO₂ from the perspective of solar energy conversion would be fundamentally interesting.

The majority of DSCs use Ru-based complexes (e.g., N₃²⁹, N₇₁₉³⁰,³¹, and black dye³²), as sensitizers of TiO₂, offering a high solar-to-electric conversion efficiency and good stability. The cost of the Ru-based sensitizers, however, is high and their synthesis and purification are complicated. Indoline-based dyes, such as D149 are promising all-organic alternatives³³. Gratzel and co-workers have produced devices based on D149, which have displayed power conversion efficiency of up to 9%³⁴. Furthermore, D149 has a peak absorption coefficient of $13900 \text{ M}^{-1} \text{ cm}^{-1}$ at 535 nm for N₇₁₉³⁵. In addition, it produces devices based on D149, which have displayed power conversion efficiency of up to 9%³⁴. Furthermore, D149 has a peak extinction co-efficient of 68700 $\text{ M}^{-1} \text{ cm}^{-1}$ at 450 nm, significantly higher than 13900 $\text{ M}^{-1} \text{ cm}^{-1}$ for N₇₁₉³⁵. In addition, it has been reported that rutile TiO₂ has a tendency for low dye loading³⁶,³⁷, even on relatively high surface area materials, compared to anatase. The use of high extinction dye should help to overcome this constraint, allowing for comparable light harvesting.

In the present paper, we report a new series of 3D hierarchical rutile TiO₂ architectures (HRT) with controlled morphologies and sized subunits by the aid of a facile acid-hydrothermal method, without any surfactant or template. We demonstrate that the hydrochloric acid (HCl) concentration plays an important role in controlling the hydrolysis rate of tetrabutyl titanate (TT) in the acid-hydrothermal process, which in turn enables us to fine-tune the final properties of rutile TiO₂. Different, well-defined 3D HRT were produced using 1–8 M HCl (denoted as HRT-1 – HRT-8, respectively). Importantly, to the best of our knowledge, this work represents the first observation of ledgewise growth for rutile TiO₂.

Furthermore, we report a solar-to-electric conversion efficiency of 5.5%, achieved using a HRT as the photoelectrode with D149 dye, which is significantly better than devices made using P25 (4.5%) and comparable to state-of-the-art commercial transparent titania anatase paste (5.8%). This power conversion efficiency is due to a considerable surface area for the high dye adsorption and a superior light scattering ability. In addition, an overall conversion efficiency 8.6% was achieved when HRT was used as the light scattering layer on top of a transparent anatase layer, a considerable improvement over commercial transparent/reflector titania anatase paste (7.6%).

Results

All the prepared samples are of the rutile phase (a = b = 0.45927 nm, c = 0.29544 nm, space group P₄₃/mmm), as shown by the XRD pattern in Supplementary Fig. S1(a), with all the peaks corresponding well with ICPDS No. 21-1276. No characteristic peaks of impurities such as anatase TiO₂, brookite TiO₂ or Ti(OH)₄ were detected and all the observed diffraction peaks suggest that the products are well crystallized. These peaks are essentially invariant with HCl concentration, as all of the morphologies appear to be of the rutile TiO₂ phase, with the Raman spectrum in Supplementary Fig. S1(b) further confirming the formation of rutile crystals. The two Raman active fundamental modes, $E_g$ ($\sim 438 \text{ cm}^{-1}$) and $A_{1g}$ ($\sim 607 \text{ cm}^{-1}$), as well as a second-order phonon as identified by Porto et al.⁶,⁷ ($\sim 247 \text{ cm}^{-1}$), and the combination of shoulders at 350, 548, and 688 cm⁻¹, which are assigned to rutile TiO₂, can be clearly distinguished.

Figure 1 (a–h) shows typical field emission scanning electron microscopy (FESEM) images of HRT samples, synthesized in a tightly sealed 45 mL Teflon-lined autoclave, from a reaction solution containing 0.5 mL precursor aqueous tetrabutyl titanate (TT) solution with 25 mL of 1 M–8 M hydrochloric acid (HCl), at 150 °C for 5 h.

For the lowest HCl concentration, 1 M, HRT-1, uniform, textured microspheres with diameters of ~1–1.5 μm are observed in low magnification SEM [Supplementary Fig. S2(a)]. It can be seen from higher magnification SEM [Fig. 1(a)] that the structure is composed of large quantities of elongated crystalline nanorods, which are closely packed and possess square profiles. FESEM of a fractured microsphere of the same material, in Supplementary Fig. S2(b), reveals that these densely packed crystalline nanorods (each with diameters of 15–30 nm) are radially arranged. It was found that the microspheres were formed by radial growth of nano-acicular crystals on the TiO₂ nanoparticulate core. Noticeably, smaller but also radially aligned nanorods form the core. It is from this surface that, nanorods with a mean width of ~20 nm and a length of ~500 nm, grew forming the TiO₂ microsphere.

The structures of HRT-2 and HRT-3 in the hydrothermal process, some-what resembling tennis balls, are 2–2.5 μm in diameter, as shown in Fig. 1(b and c). The domain like arrangements of the TiO₂ nanorods form quasi spheres leading to an apparently smooth
surface. The surfaces of these spheres are in fact densely packed nanorods. Samples prepared with 4 M HCl resulted in the formation of morphology likened to a cauliflower, comprised of the near-perfect assemble of numerous TiO$_2$ nanorods, as the nanorod bundles are smaller and more disparate [Fig. 1(d)].

SEM of samples HRT-5 and HRT-6 [Fig. 1(e and f)] revealed even more open morphologies with the nanorod bundles more separated from each other. Each of these bundles is assembled from several dozen nanorods of lengths of up to 3 μm. The distribution of quasi-cubic primary nanocrystallites with sizes 15–25 nm led to the observation of surfaces which may be likened to flowers. These differences are due to the initial differences in crystal growth of the titanium precursor in the stronger acidic medium (described in more detail below). The observation of nanostructures indicate that the formation of prisms and the flower-like structures, possibly stemmed from the crystal growth of the primary nanoparticles and the further epitaxial process of these prisms.

High HCl concentrations of HRT-7 and HRT-8 resulted in even more open structures again, as the surface etching rate is much faster and thus tends to result in more nucleation sites at early stages of synthesis. The resulting nanorods bundles are thinner as they contain fewer rods [Fig. 1(g, h)]. At these concentrations, structures composed of hundreds of nanoneedles were produced. It is known that high HCl concentration is favorable for the formation of smaller-scale building blocks.

To elucidate the microstructural origin of the growth mechanisms, the architectures of the HRT obtained under different HCl concentrations were further investigated by high resolution SEM, high resolution transmission electron microscopy (HRTEM) and fast Fourier transforms (FFTs) of the HRTEM images. Figure 2(a) presents a high-magnification SEM image obtained from the synthesis completed using 0.5 mL of TT and 1 M HCl. Figure 2(b) shows the TEM image of a single nanorod (prepared under the same conditions) with a width of 20 nm. The FFT pattern is shown in inset of Fig. 2(b) and, along with the associate HRTEM image, confirm the single-crystal nature of the rutile TiO$_2$ rods.

It can be observed from Fig. 2 that the rutile crystals (HRT-1 to HRT-5) can be described by a combination of an elongated cuboid structure and pyramid shaped cap [this is also seen in Fig. 3(a) and schematically in Supplementary Fig. S3(a)]. Two sets of lattice fringes perpendicular to each other with distances of 0.29 and 0.32 nm can be readily ascribed to the lattice spacings of the (001) and (110)
planes of the rutile TiO₂, respectively [Supplementary Fig. S3(b)]. The FFT pattern [inset of Fig. 2(b)] further confirms that (1) the cuboid crystal facets are parallel to [110], (2) the pyramid-shaped crystal facets are parallel to [111], and (3) a preferred growth along the [001] direction takes place. Moreover, according to periodic bond chain (PBC) theory, PBCs result in the formation of the crystal and control the direction of crystal growth (the strongest chemical bonds are usually the ones at which the crystal has the fastest growth rate). Theoretical calculation shows that for a rutile lattice, the surface energy is \( E(110) < E(100) < E(101) < E(001) \). In crystal growth, a fast growing plane generally tends to disappear leaving behind slower growing planes with lower surface energies. Thus, crystal growth along the [001] direction (i.e., \( c\)-axis) would have the fastest growth rate compared to other directions.

Interestingly, the top of all the rods are pyramidally capped, as shown in Fig. 2. HRTEM imagery of the rod [Fig. 2(b)] indicates that the plane of the pyramid has a contact angle of approximately 130° to the lateral [110] plane. The most appropriate plane is considered to be [111] under this condition because contact angles between [110] and [101], [110] and [111], and [110] and [221] planes are calculated to be 112.5°, 132.3°, and 151.2°, respectively, using the tetragonal rutile unit cell.

No crystallographic differences were observed between HRT-1 and HRT-8. Our results then demonstrate that [111] and [110] faceted rods are obtained during the hydrothermal growth. It should be noted that the [111] planes observed at the top surface of the rods are a minor surface in the equilibrium shape of a macroscopic rutile TiO₂ crystal using the Wulff construction and the calculated surface area is a minor surface in the equilibrium shape of a macroscopic rutile.

It has been demonstrated previously that Cl⁻ plays an important role in the growth of rutile TiO₂ nanorods. This is because the [110] planes of rutile have a positive polar face and Cl⁻ is preferentially adsorbed on this surface, which serves to restrain the contact of TiO₂ grains on the [110] surfaces and thus greatly retards crystal growth along the [001] plane. As a result, TiO₂ grain growth was suppressed in the \(<110>\) directions and accelerated in the [001] direction, thereby forming rod-like structures. However, our results have shown the presence of the [111] facet with a large surface area, giving rise to the pyramid-type morphology. The authors believe these results give important new insight into the surface chemistry of rutile TiO₂. It is reasonable to expect that the TiO₂ nanorods preferentially exposes the [110] side facets and the [111] top facets and grows along the [001] direction on the basis of the EPR and HRTEM results together with the XRD pattern.

Figure 2 (c–j) shows the morphology and structure of the nanorods bundles and individual rods for HRT-2 to HRT-5. To further obtain structural information for the radial rods, TEM images and the corresponding FFT patterns are also recorded on a single nanorod of each (for HRT-6 to HRT-8, however, only bundles of nanorods were taken because the samples were brittle and no loose bundles could be identified, probably because of bundles more open but rods within each bundle more tightly held together). Based on Fig. 2 it can be seen that, similar to HRT-1, the TiO₂ nanorod synthesized with 2 M–5 M HCl, also preferentially exposes the [110] side facets and the [111] top facets and grows along the [001] direction. In addition, the nanorod crystals of the HRT synthesized in the highest HCl concentration solutions (6 M–8 M) grew along the [001] direction. The exposed surfaces of the nanorods are [110] facets. The variation of HCl concentration in the reaction solutions did not affect the crystal growth direction of the nanorods.

In spite of the differences in micron-scale morphologies, all the architectures are composed of very similar crystalline nanorod units. It has been demonstrated previously that Cl⁻ plays an important role in the growth of TiO₂ nanorods (as opposed to nanoparticles). This is because the [110] planes of rutile have a positive polar face and Cl⁻ is preferentially adsorbed on this surface, which serves to restrain the contact of TiO₂ grains on the [110] surfaces and thus greatly retards crystal growth along the [001] plane. As a result, TiO₂ grain growth was suppressed in the \(<110>\) directions and accelerated in the [001] direction, thereby forming rod-like struc-

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**Figure 3 | Ledgewise growth of rutile nanorod (HRT-1).** (a) Macro growth direction of a rutile nanorode; preferential growth to [001] direction, (b) Thickening of a nanorode by ledgewise growth of (110) surface, (c) atomic structure of a rutile nanorod.
It is interesting that a ledgewise growth mechanism was observed for HRT-1, as shown in Fig. 3. Preferential growth of a rutile nanorod to the [001] direction is also confirmed by overview of a nanorod [Fig. 3(a)]. The growth rate in the [001] direction is much faster than that of the [110] direction, resulting in an elongated morphology in the [001] direction. Lateral growth of a nanorod, i.e., the diameter increasing, occurs with the generation of a fresh [110] surface. High resolution images of the side of a nanorod show several steps on a (110) plane [Fig. 3(b), Supplementary Fig. S3(b)]. Importantly, these steps have a step height of 0.32 nm, coincident with interplanar spacing of a (110) plane. This ledgewise growth mechanism minimizes the total surface energy during growth of nanorod [Fig. 3(c)], as a monolayer step will minimize the amount of high surface energy area. The total (110) surface area which has minimum surface energy is kept constant during growth. The growth of a rutile nanorod is displayed schematically in Supplementary Fig. S3(a). The cuboidal body has [110] type surface and the upper pyramidal cap is composed of [111] planes.

Brunauer-Emmett-Teller (BET) calculations of specific surface area based on N$_2$ adsorption-desorption were performed (Supplementary Table S1), however, HRT-5–HRT-8 were omitted from this due to the low yields which lead to large measurement errors. The specific surface area is found to decrease with increasing concentration of HCl used in synthesis. N$_2$ adsorption-desorption isotherms (shown in Supplementary Fig. S4) of the samples (HRT-1 – HRT-4) exhibited typical type IV sorption behavior with hysteresis loops, ascribed to the existence of the mesoporous structure.

On account of it having the highest specific surface area (and highest yield among all the samples), HRT-1 was applied in photoanodes for DSCs. Two control materials were employed for comparative purposes. Firstly, P25 was used, as it is low cost and widely reported in this field. P25 also represents a mixed phase material (mostly anatase, with a significant rutile impurity). A commercial transparent titania anatase paste [18NR-T, denoted as Dyesol-T, obtained from Dyesol (Australia)] was also chosen for its role in some of the highest reported DSC efficiencies to date.

The porosity (P) and surface roughness factor (R) based on adsorption-desorption measurements were measured for the two controls and compared with HRT-1, summarized in Table 1. HRT-1 is >50% porous, which is promising as porosities of nanocrystalline TiO$_2$ films employed in high efficiency DSCs are typically in the range of 50–65%. HRT-1 has a similar specific surface area as Dyesol-T (74 m$^2$ g$^{-1}$), but significantly higher than P25. Furthermore, the calculated roughness factor for HRT-1, approximately 125 µm$^{-1}$, which is higher than both P25 and Dyesol-T.

The photocurrent density-voltage (J-V) curves of P25, HRT-1 and Dyesol-T based DSCs (each with ~12 µm porous TiO$_2$ layer), are shown in Fig. 4, (summarized in Supplementary Table S2). From this, it can be found that the overall energy conversion efficiency (η) of HRT-1 was higher than P25 but lower than Dyesol-T. The short-circuit photocurrent density ($J_{sc}$) is highest in the devices with Dyesol-T, then HRT-1, and lowest when P25 is used. This order is repeated in the dye loading data collected (Table 1, Supplementary Fig. S6), however is probably not the whole story with regard to light harvesting efficiency (LHE). The higher $J_{sc}$ for HRT-1 based devices than P25 could be mainly attributed to the improved dye loading capacity, attributable to measured R and BET surface area results, as well as light scattering properties of the TiO$_2$ films. In both cases the materials are scattering. It can be expected that this will actually cause some problems, with reflection near to the front surface causing light to be lost.

Light scattering generally serves to improve LHE. As shown in Fig. 4(b), HRT-1 is a much more effective light scattering material than either P25 or Dyesol-T (the reflectance spectrum from 400 to 800 nm). In order to better scrutinize the different photovoltaic performances of these three solar cells, the incident photon-to-current efficiency (IPCE) spectra as a function of wavelength was collected. The IPCE [inset of Fig. 4(a)] is defined as the ratio of the number of electrons extracted the external circuit (under short circuit conditions) divided by the number of incident photons as a function of the wavelength of incident light. As such, it is a function of LHE as well as charge injection, $\eta_{inj}$ and charge transport efficiencies, $\eta_{tr}$.

It is observed that devices made with HRT-1 possess a marginally lower peak IPCE values than that of Dyesol-T devices and significantly higher than P25. This is an important point to observe as it suggests that $\eta_{inj}$ for HRT-1 should be similar to Dyesol-T (and close to unity, given the light transmission of the fluorine-doped tin oxide (FTO) substrate). The lower IPCE values recorded away from the dye peak absorbance wavelength can be attributed to either lower $\eta_{tr}$ (addressed below) or light being reflected out of the cell before it can be absorbed by a dye molecule. At the maximum value of the IPCE spectrum at around 530 nm, the IPCE of the HRT-1 film is approximately 10% higher than that of the P25 film. This increase is in good agreement with the increased $J_{sc}$ and $\eta$ observed in Fig. 4(a) and Supplementary Table S2.

Measured $V_{oc}$ under one sun illumination follows the same trend as $J_{sc}$ (Dyesol-T > HRT-1 > P25). Having said this, the difference between the pure anatase and pure rutile devices here is much less (~30 mV) than would be expected based solely on conduction band offsets (~200 mV). Perhaps the most surprising aspect of these results however is the enhanced fill factor (FF) for HRT-1 compared to the other two materials. This is analyzed in greater detail below.

The highly faceted nature of these structures results in less defect sites compared to flame pyrolysed P25, which should serve to decrease electron recombination for TiO$_2$ to I$_3^-$ in the electrolyte. Electrochemical impedance spectroscopy (EIS) is a very useful tool for understanding the kinetics of electrochemical and photoelectrochemical processes occurring in DSCs$^{4-8}$. Fig. 4(c) presents the Nyquist plots of the impedance spectra obtained under illumination for DSCs assembled with P25, HRT-1 and Dyesol-T films under $V_{oc}$. and the equivalent circuit is shown as the inset. The figure shows Nyquist plots with two semicircles. All R values are similar (7.0–7.5 Ω). The left semicircle (high frequency) represents the resistance at the counter electrode ($R_{ct}$) for the reduction reaction of I$_3^-$ ions in the electrolyte using Pt as the counter electrode. The second semi-

| Table 1 | Specific surface area, porosity, and roughness factor of P25, HRT-1 and Dyesol-T materials |

| Samples | Specific surface area (m$^2$ g$^{-1}$) | Porosity$^{a}$ (%) | Roughness factor$^{b}$ (µm$^{-1}$) | Amount of dye$^{c}$ (10$^{-7}$ mol cm$^{-2}$) |
|---------|-------------------------------------|------------------|---------------------------------|-----------------------------------|
| P25     | 42                                 | 34.9             | 118                             | 2.4                              |
| HRT-1   | 67                                 | 56.6             | 125                             | 2.7                              |
| Dyesol-T| 74                                 | 64.2             | 103                             | 3.4                              |

$^{a}$The porosities (P) of P25, HRT-1 and Dyesol-T were calculated according to: $P = \frac{V_p}{V_i + V_e}$, where $V_p$ is the specific cumulative pore volume (cm$^3$ g$^{-1}$) and $V_i$ is the density of TiO$_2$ ($\rho = 4.3$ g cm$^{-3}$).$^{10}$

$^{b}$An estimation of the roughness factor (R) per unit film thickness of the films is obtained by $R = \mu (P - 1) - P S$, where $\mu$ is the density (g cm$^{-3}$) of TiO$_2$, $P$ is the porosity (%) of the film, and $S$ is the specific surface area (m$^2$ g$^{-1}$).$^{11}$

$^{c}$The dyed electrodes were soaked in a 0.4 M alkaline solution in methanol to desorb the dye from the electrodes. The amount of desorbed dye was quantified by measuring its optical absorption spectrum (Supplementary Fig. S5d).
Figure 4 | (a) J-V characteristics of P25-, HRT-1- and Dyesol-T-based DSCs. Cells were illuminated at an intensity of 100 mW cm\(^{-2}\) with a spectrum approximately AM 1.5 G and an active area of 0.16 cm\(^2\); inset: incident photon to current conversion efficiency (IPCE) curves of P25-, HRT-1- and Dyesol-T-based DSCs. (b) diffuse reflectance spectra of P25, HRT-1 and Dyesol-T films. Impedance spectra of DSCs containing P25, HRT-1 and Dyesol-T photoanodes measured at \(V_{oc}\) under illumination of 100 mW cm\(^{-2}\); (c) Nyquist plots, with the inset showing the equivalent circuit, and (d) Bode phase plots.

circle in the intermediate frequency range results from charge transport through the semiconductor, while the corresponding chemical capacitance \(C_D\) describes the density of states in the bandgap of TiO\(_2\). These parameters were calculated from measured data by using an equivalent circuit [inset of Fig. 4(c)] as shown in Supplementary Fig. S7 and Supplementary Table S3.

According to Adachi et al., several parameters related to the properties of electron transport in the semiconductor can be deduced from the Nyquist plot. In particular, the charge transport resistance \(R_{ct}\) related to electron-hole recombination can be determined from the central arc diameter. In our case, it was found that \(R_{ct}\) at the TiO\(_2\)/dye/electrolyte interface for HRT-1 was much smaller than for P25 and Dyesol-T, implying a relatively faster electron transfer process and a slower electron recombination at the HRT-1 photoanode interface. Meanwhile, the higher \(C_D\) value of the HRT-1-based DSC indicates that more states on the TiO\(_2\) surface are able to accept electrons from the excited state of the dye, yielding an increased photocurrent. Furthermore, the Bode phase plots shown in Fig. 4(d) likewise support the differences in the electron lifetime for the TiO\(_2\) films derived from P25, HRT-1 and Dyesol-T. The electron lifetime \(\tau_{eff}\) could be calculated from the low frequency results as:

\[
\tau_{eff} = \frac{1}{2\pi f_{max}}
\]

where the characteristic \(f_{max}\) is the maximum phase shift in the mid-frequency peak. The middle-frequency peaks of the DSCs containing P25 and Dyesol-T shift to higher frequency relative to HRT-1, indicating a longer electron lifetime for the latter material. The longer electron lifetime observed with HRT-1 sensitized solar cells indicates more effective suppression of the back-reaction of the injected electrons with the \(I_3^-\) ions in the electrolyte. With these considerations, it seems likely that the limitations of the rutile materials are not connected to \(\Phi_{oc}\), implying that reflecting light out of the cell before it can be absorbed is the major limiting factor here. It is quite interesting to note that charge transport through HRT-1 is better than for the other materials, in spite of the bulk charge transport values being lower in rutile than anatase. The longer electron lifetime is likely due to morphology of rods in the larger constructs.

Previous studies showed that, lower dye loadings (N719) on rutile compared to anatase, for the same thickness, is the main reason for lower \(f_{ac}\) and \(\eta\). However, using organic sensitizer (D149) in our study allows us to realize similar light harvesting for both rutile and anatase.

As an additional experiment, the light scattering properties of HRT-1 were compared to a commercially available reflector titania anatase paste [WER2-O, denoted as Dyesol-S, obtained from Dyesol (Australia)]. Once again, this has been used by a number of groups around the world in order to produce high efficiency devices.

Both scattering layers were applied on top of a \(\sim 12\) \(\mu\)m thick layer of Dyosol-T, with diffuse reflectance spectra of Dyosol-T (\(\sim 12\) \(\mu\)m)/Dyesol-S (\(\sim 4\) \(\mu\)m) and Dyesol-T (\(\sim 12\) \(\mu\)m)/HRT-1 (\(\sim 4\) \(\mu\)m) films shown in Supplementary Fig. S8. Both these materials provided enhanced performance efficiencies, with the most marked improvement being seen when HRT-1 is employed. Figure 5 (summarized in Supplementary Table S4) shows the J-V curves of DSCs using the two scattering layers. A \(J_{sc}\) of 15.5 mA cm\(^{-2}\), \(V_{oc}\) of 740 mV, \(FF\) of 75.0%, and an overall \(\eta\) of 8.6% was obtained for Dyesol-T/HRT-1 films, under AM 1.5 G one sun intensity, which is superior than the optimized Dyosol-T/Dyesol-S (7.6%). Much of this improvement can be attributed to higher \(FF\) when HRT-1 is used, compared to Dyesol-S. Investigations into the root cause of this are ongoing.

Discussion

In summary, different well-defined 3D hierarchical rutile TiO\(_2\) architectures (HRT) have been synthesized via a facile hydrothermal method without any surfactant or template. A ledgewise growth
was observed for the first time on rutile TiO2. A 3D architecture has been demonstrated as an effective photocathode when employed in a DSC along with a metal-free organic dye (D149). Although the performance of HRT based devices was inferior to those of state-of-the-art commercial anatase Dyesol-T paste, this study provided important insights into rutile TiO2 and, in particular its potential applications in DSCs. The difference in Jsc here is significantly less than in previous studies25. Charge transport is seen to be less of a problem than it was initially anticipated to be. Further investigations and comparison between rutile and anatase are ongoing, including looking into bi-layer rutile devices incorporating a less scattering front layer.

An energy conversion efficiency of 5.5% was achieved, which is superior to that of an analogous device using P25 (4.5%), and comparable to Dyesol-S (5.8%). The overall conversion efficiency 8.6% was achieved when HRT was used as the light scattering layer, a considerable improvement over commercial anatase Dyesol-T (~12 μm)/Dyesol-S (~4 μm) system (7.6%). The 3D rutile TiO2 and metal-free indoline sensitizer may present an alternative candidate with regard to producing high-performance DSCs.

**Methods**

**Preparation of 3D hierarchical rutile TiO2 architectures (HRT).** Tetrabutyl titanate (Ti(OCH2CH2CH2CH3)4, 97%, T3, analytical reagent grade) and hydrochloric acid (37% HCL, analytical reagent grade) were purchased from Sigma-Aldrich, and used without purification or additives. The typical experimental procedures are described as follows: firstly, 0.5 mL of TT was added dropwise into an amount of concentrated HCl under magnetic stirring. Then, the upper solution was diluted by distilled water to meet a total volume of 25.5 mL under different HCl concentrations (1 M–8 M). After vigorously stirring for another 1.5 h and aging for 1 h, the whole mixture was transferred into a 45 mL Teflon-lined stainless steel autoclave and sealed. The hydrothermal synthesis was conducted at 150 °C for 30 min, and then at 500 °C for 15 min) to remove organic components and improve interparticle connectivity, the TiO2 films were soaked in the dye bath, washed with acetic acid, and dried. The working electrode and Pt counter electrode [produced using a pre-dried piece of 2.5 mm FTO glass, coated with one drop of 10 m platonic acid solution [H2PtCl6 (Sigma)] and heated to 400 °C for 20 min] were assembled into a sandwich type cell and sealed with a spacer of 25 μm Surlon (Solaronix). At 1.5 V reverse bias, the I-V curves were measured using a Keithley 2400 sourcemeter. The measured currents were referenced to a calibrated Si solar cell. The short circuit current of the devices was recorded in 5 nm steps using a Keithley 2400 sourcemeter. The measured currents were calibrated to a reference current density of a standard Si photodiode (PECCELL).

**Materials characterizations.** The crystal phase of the 3D hierarchical TiO2 architectures were measured with a powder X-ray diffractometer (XRD, MMA, GBC Scientific Equipment LLC, Hampshire, IL, USA) with Cu Kα radiation. The morphology of the samples was examined by Field Emission Scanning Electron Microscopy (FESEM, JSM-7500FA, JEOL, Tokyo, Japan) and High Resolution Transmission Electron microscopy (HRTEM, JEM-2100F). Brunauer–Emmett–Teller (BET) surface area (SBET) were determined by using a nitrogen adsorption apparatus (Tristar 3030, Micromeritics Instrument Corporation), with all samples being degassed at 120 °C overnight before measurement. Dy dispersion was performed by dipping the sensitized sample into 0.4 M NaOH in methanol solution. Film thickness was measured with a Veeco Dektak 150 Surface Profiler. The light scattering properties of the samples were investigated by ultraviolet (UV)–visible light absorption/diffuse reflectance spectrometry (Shimadzu UV-3600) with integrating sphere attachment.

**Preparation of TiO2 photoanodes.** To prepare the DSC working electrodes, fluorine doped tin oxide (FTO) glass (2.3 mm thick, 8 Ω/sq, Hartford Glass), was used as current collector, was first cleaned with soapy water for 20 min, rinsed with distilled water, immersed in an ultrasonic bath in acetone and then again in ethanol for 30 min each. Following this, a dense TiO2 blocking layer was deposited on the clean FTO glass by spray pyrolysis of a solution of diluted titanium diisopropanoxide bis(acetylacetonate), 75% in isopropanol (Aldrich), in absolute ethanol in the volume ratio of 1:9. Next, a layer of TiO2 paste was cast onto the FTO glass plates by a manual doctor–blade method. These pastes were prepared as follows: 0.1 g of TiO2 powder was ground with a mixture of ethanol (10 mL), distilled water (10 mL), nitric acid (0.167 mL), polyethylene glycol solution (4 mL, 10 wt% in distilled water, MW = 100,000), and triton X-100 (one small drop) to form a slurry, and then the slurry was sonicated for 1 h in an ultrasonic bath and stirred for 2 h. A viscous white TiO2 paste was finally obtained. After a heating process (at 150 °C for 10 min, at 325 °C for 5 min, at 375 °C for 10 min, at 450 °C for 30 min, and then at 500 °C for 15 min) to remove organic constituents and improve interparticle connectivity, the TiO2 films were soaked in 0.02 M aqueous titanium tetrachloride solution (TiCl4) at 70 °C for 30 min, before being sintered again (500 °C for 30 min). State-of-the-art commercial transparent titania anatase paste (18NR-T, denoted as Dyesol-T), and commercial reflector titania anatase paste (WBR2-O, denoted as Dyesol-S), [both obtained from Dyesol (Australia)] were used for the references.

**Fabrication of Dye-sensitized solar cells.** The porous TiO2 films were immersed in a 0.5 mM D149 (1-material, Canada) dye solution in a 1:1 (v/v) mixture of acetonitrile (HPLC, Lab-scan) and tert-butanol (LR, Ajax Chemicals) overnight once their temperature decreased to approximately 110 °C. The samples were then taken out of the dye bath, washed with acetic acid, and dried. The working electrode and Pt counter electrode [produced using a pre-dried piece of 2.5 mm FTO glass, coated with one drop of 10 m platonic acid solution [H2PtCl6 (Sigma)] and heated to 400 °C for 20 min] were assembled into a sandwich type cell and sealed with a spacer of 25 μm Surlon (Solaronix). At 1.5 V reverse bias, the I-V curves were measured using a Keithley 2400 sourcemeter, a simulated 100 mW cm−2 air mass (AM) 1.5 light source (Oriel), and customized LabView software. The device area was masked with a metallic mask slightly larger than the active area. Incident photon-to-current quantum conversion efficiency (IPCE) was measured using a 300 W Xe lamp, a monochromator with sorting filters, focused to a spot with additional optics. The short circuit current response of the devices was recorded in 5 nm steps using a Keithley 2400 sourcemeter. The measured currents were calibrated to a reference current density of a standard Si photodiode (PECCELL).

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Author contributions
J.L., A.N. and J.H.K. designed experiments. J.L. performed synthesis experiments and DSC measurements. Y.U.H. performed TEM. Z. S. supervised the synthesis experiment. J.L., A.N., L.W. and J.H.K. analysed the data. J.L., A.N., J.H.K. and S.X.D. wrote the manuscript. All authors discussed the results and contributed to revisions.

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