Short-Term Synaptic Dynamics Control the Activity Phase of Neurons in an Oscillatory Network

Abbreviated Title: Synaptic dynamics influence activity phase of oscillatory neurons

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Abstract

In oscillatory systems, neuronal activity phase is often independent of network frequency. Such phase maintenance requires adjustment of synaptic input with network frequency, a relationship that we explored using the crab, *Cancer borealis*, pyloric network. The burst phase of pyloric neurons is relatively constant despite a >2-fold variation in network frequency. We used noise input to characterize how input shape influences burst delay of a pyloric neuron, and then used dynamic clamp to examine how burst phase depends on the period, amplitude, duration, and shape of rhythmic synaptic input. Phase constancy across a range of periods required a proportional increase of synaptic duration with period. However, phase maintenance was also promoted by an increase of amplitude and peak phase of synaptic input with period. Mathematical analysis shows how short-term synaptic plasticity can coordinately change amplitude and peak phase to maximize the range of periods over which phase constancy is achieved.

Introduction

Oscillatory neural activity is often organized into different phases across groups of neurons, both in brain rhythms associated with cognitive tasks or behavioral states (Hasselmo et al., 2002; Buzsaki and Wang, 2012; Buzsaki and Tingley, 2018), and in central pattern generating (CPG) circuits that drive rhythmic motor behaviors (Marder and Bucher, 2001; Marder et al., 2005; Grillner, 2006; Bucher et al., 2015; Katz, 2016; Stein, 2018). The functional significance of different phases in the latter is readily apparent, as they for example provide alternating flexion and extension of limb joints, and coordination of movements between joints, limbs, and segments (Krantz and Parks, 2012; Grillner and El Manira, 2015; Kiehn, 2016; Le Gal et al., 2017; Bidaye et al., 2018). A hallmark of many such patterns is that the relative timing of firing between neurons is well maintained over a range of rhythm frequencies (Dicaprio et al., 1997; Hooper, 1997b, a; Wenning et al., 2004; Marder et al., 2005; Grillner, 2006; Mullins et al., 2011; Le Gal et al., 2017). If the latency of firing across different groups of neurons changes proportionally to the rhythm period, phase (latency over period) is invariant, in some cases providing optimal limb coordination at all speeds (Zhang et al., 2014).
The ability of the system to coordinate phases with changes in period arises from central coordinating mechanisms between circuit elements, as it is present in isolated nervous system preparations, but the underlying cellular and circuit mechanisms are not well understood. For instance, constant phase lags between neighboring segments in the control of swimming in lamprey fish and crayfish can be explained mathematically on the basis of asymmetrically weakly coupled oscillators, but the role of intrinsic and synaptic dynamics within each segment is unknown (Cohen et al., 1992; Skinner and Mulloney, 1998; Grillner, 2006; Mullins et al., 2011; Zhang et al., 2014; Le Gal et al., 2017).

The pyloric circuit of the crustacean stomatogastric ganglion (STG) has inspired a series of experimental and theoretical studies of cellular and synaptic mechanisms underlying phase maintenance. The pyloric circuit generates a triphasic motor pattern with stable phase relationships over a wide range of periods (Eisen and Marder, 1984; Hooper, 1997b, a; Bucher et al., 2005; Goaillard et al., 2009; Tang et al., 2012; Soofi et al., 2014). Synapses in the pyloric circuit use graded as well as spike-mediated transmission (Graubard et al., 1980; Harris-Warrick and Johnson, 2010; Zhao et al., 2011; Rosenbaum and Marder, 2018). Follower neurons burst in rebound from inhibition from pacemaker neurons (Marder and Bucher, 2007; Daur et al., 2016), and post-inhibitory rebound delay scales with the period of hyperpolarizing currents (Hooper, 1998). Voltage-gated conductances slow enough for cumulative activation across cycles could promote such phase maintenance (Hooper et al., 2009). Similarly, short-term depression of graded inhibitory synapses is slow enough to accumulate over several pyloric cycles, meaning that effective synaptic strength increases with increasing cycle period (Manor et al., 1997; Nadim and Manor, 2000).

Theoretical studies have shown that short-term synaptic depression, by increasing inhibition strength with cycle period, should promote phase maintenance (Manor et al., 2003; Mouser et al., 2008), particularly in conjunction with inactivating (A-type) potassium currents (Bose et al., 2004; Greenberg and Manor, 2005), which control the rebound delay (Harris-Warrick et al., 1995b; Harris-Warrick et al., 1995a; Kloppenburg et al., 1999). These predictions remain experimentally untested.

Additionally, postsynaptic responses also depend on the actual trajectory of synaptic conductances, which are shaped by presynaptic voltage trajectories and short-term synaptic plasticity (Manor et al., 1997; Mamiya et al., 2003; Zhao et al.,
If amplitude, duration, and trajectory of synaptic conductance determine rebound delay, phase maintenance necessitates all three of these parameters to change with cycle period in coordination. We used the dynamic clamp technique to exhaustively explore the range of these parameters and understand how the coordinated changes in synaptic dynamics determines the phase of follower neurons in an oscillatory circuit. Our findings are consistent with a mathematical framework that accounts for the dependence of amplitude and peak phase of the synaptic conductance on cycle period.

**Results**

**Phase maintenance and latency maintenance**

The firing of neurons in oscillatory networks is shaped by a periodic synaptic input. The relative firing latency of such neurons is often measured relative to a defined reference time in each cycle of oscillation, and is used to determine the activity phase of the neuron (see, e.g., Belluscio et al., 2012). For example, in a simple network consisting of a bursting oscillatory neuron driving a follower neuron (Fig. 1A1), at a descriptive level, the latency ($\Delta t$) of the follower neuron activity relative to the onset of the oscillator’s burst onset may depend on the oscillation cycle period ($P$). In response to a change in period (say, to $P_2$), the follower neuron may keep constant latency ($\Delta t_2 = \Delta t$), or constant phase, i.e., modify its latency proportionally to the change in period ($\Delta t_2 / P_2 = \Delta t / P$; Fig. 1A2). However, in many oscillatory systems, for example the pyloric circuit (Hooper, 1997b, a), the relationship between $L$ and $P$ falls between these two extremes.

We demonstrated this point in the pyloric follower LP neuron using the following protocol. We voltage clamped one of the pacemaker PD neurons and drove this neuron with its own pre-recorded waveform, but applied at five different cycle periods ($P$). This protocol entrained the pacemaker group at this period, which forced the follower LP neuron to obey the same period (Fig. 1B). We then measured the latency ($\Delta t$) of the LP burst onset with respect to onset of the PD neuron burst. A plot of the LP latency $\Delta t$ or phase ($\Delta t / P$) for different cycle periods demonstrates the above-mentioned finding that the LP neuron activity falls between the two limits of constant phase and constant latency (Fig. 1C).
The burst onset time of the LP neuron depends on the temporal dynamics of its input

The LP neuron does not have intrinsic oscillatory properties, but oscillates due to the synaptic input it receives from the pacemaker anterior burster (AB) and pyloric dilator (PD) neurons, and the follower pyloric constrictor (PY) neurons (Fig. 2A). The burst onset phase of the LP neuron \( \phi_{LP} = \frac{\Delta t}{P} \) (Fig. 2A) is shaped by the interaction between synaptic inputs and the neuron’s intrinsic dynamics that influence post-inhibitory rebound. We measured an overall burst onset phase of the LP neuron of \( \phi_{LP} = 0.34 \pm 0.03 \) (N=9).

As a first-order quantification, we measured how inputs to the LP neuron interact with its intrinsic properties to determine the timing between its bursts, in the absence of network oscillations. To this end, we blocked the synaptic input from the pacemaker AB and follower PY neurons to the LP neuron (Fig. 2B) and drove the LP neuron with a noise current input (see Methods). In response to the noise input, the LP neuron produced an irregular pattern of spike times, which included a variety of bursting patterns with different spike numbers (Fig. 2C). We were interested in the characteristics of inputs producing different burst onset latencies. However, unlike a periodic input, noise input does not provide a well-defined reference point to measure the burst onset latency. We categorized bursts with respect to the preceding inter-burst intervals (IBIs; see Methods), during which no other action potentials occurred. We classified these IBIs in bins (300, 500, 700 and 900 ms) and tagged bursts based on the IBI values (Fig. 2C). We characterized the driving input leading to bursts with specific IBIs by burst-triggered averaging the input current \( I_{BTA} \); an example shown in Fig. 2D). Our analysis produced a single \( I_{BTA} \) for each of the four IBIs in each preparation (N=23). \( I_{BTA} \)'s of each preparation were first normalized in amplitude by the (negative) peak value of the \( I_{BTA} \) at IBI = 300 ms (Fig. 2E; average shown in Fig. 2F) to examine how peak amplitude \( I_{peak} \) varied with IBI. These data were then normalized in time (Fig. 2G) to examine the effect of IBI on peak phase \( \Delta_{peak} \) and the rise \( slope_{up} \) and fall \( slope_{down} \) slopes of the input current across preparations. We found that IBI had a significant effect on \( I_{peak} \), \( \Delta_{peak} \), \( slope_{up} \) and \( slope_{down} \) (all one-way RM-ANOVA on ranks; data included in Figure 2-source data). In particular, larger IBIs corresponded to larger \( I_{peak} \) values (Fig. 2F-2H; p<0.001, \( \chi^2 = 65.87 \) with smaller (more advanced) \( \Delta_{peak} \) (Fig. 2I; p<0.001, \( \chi^2 = 41.35 \)). The change in \( \Delta_{peak} \) was due to a decrease in
The burst onset phase of the LP neuron oscillation depends on its synaptic input

Injection of noise current revealed that the timing of the LP response is exquisitely sensitive to the duration and amplitude of inputs. In the intact system, the primary determinant of input duration and amplitude is the network period \( P \), as increasing \( P \) increases both presynaptic pacemaker burst duration (Hooper, 1997b, a) and synaptic strength (Manor et al., 1997; Nadim and Manor, 2000). To explore the effect of the duration and strength of the synaptic input, we used dynamic clamp to drive the LP neuron with a realistic synaptic conductance waveform.

We constructed this realistic waveform by measuring the synaptic current input to the LP neuron during ongoing pyloric oscillations (Fig. 3A). These measurements showed the two components of inhibitory synaptic input: those from the pacemaker AB and PD neurons (left arrow) and those from the follower PY neurons (right arrow). In each cycle, the synaptic current always had a single peak, but the amplitude and phase of this peak showed variability across preparations (Fig. 3B, average in blue).

The realistic conductance input was injected periodically with strength \( g_{\text{max}} \) (Fig. 3C). For any fixed \( g_{\text{max}} \), \( \phi_{\text{LP}} \) decreased as a function of \( P \) (Fig. 3D), i.e., the relative onset of the LP burst was advanced in slower rhythms. In contrast to the effect of \( P \), for any given \( P \), \( \phi_{\text{LP}} \) increased sublinearly as a function of \( g_{\text{max}} \) (Fig. 3E). Fig. 3F combines the simultaneous influence of both parameters on \( \phi_{\text{LP}} \). The results shown in Fig. 3D indicate that the LP neuron intrinsic properties alone do not produce phase constancy. However, level sets of \( \phi_{\text{LP}} \) (highlighted for three values in Fig. 3F), indicate that phase could be maintained over a range of \( P \) values, if \( g_{\text{max}} \) increases as a function of \( P \). This finding was predicted by our previous modeling work, in which we suggested that short-term synaptic depression promotes phase constancy by increasing synaptic strength as a function of \( P \) (Manor et al., 2003; Bose et al., 2004). We will further discuss the role of synaptic depression below.

To clarify the results of Fig. 3, it is worth examining the extent of phase maintenance for fixed \( g_{\text{max}} \). An example of this is shown in Fig. 4A (turquoise plots). A comparison of these data with the theoretical cases in which either delay or phase is constant suggests that the LP neuron produces relatively good phase maintenance, at
least much better in comparison with constant delay. However, this conclusion is misleading because, in these experiments, the duty cycle of the synaptic input was kept constant. Therefore, most of the phase maintenance is due the fact that the synaptic input keeps perfect phase. In fact, if the reference point measures phase relative to the end—rather than onset—of the PD burst (Fig. 4B), phase maintenance of the LP neuron is barely better than in the constant delay case (Fig. 4A, purple plots). It is therefore clear that phase maintenance by the LP neuron would require the properties of the synaptic input to change as a function of $P$, a hallmark of short-term synaptic plasticity (Fortune and Rose, 2001; Grande and Spain, 2005). As mentioned above, short-term plasticity such as depression could produce changes in $g_{\text{max}}$ as a function of $P$. Independently of $g_{\text{max}}$, the peak time of the synaptic current is another parameter that could change with $P$ and influence the timing of the postsynaptic burst. We therefore proceeded to systematically explore the influence of $P$, $g_{\text{max}}$ and the synaptic peak time on $\phi_{\text{LP}}$.

**A systematic exploration of synaptic input parameters on the phase of the LP neuron**

For a detailed exploration of the influence of the synaptic input on $\phi_{\text{LP}}$, we approximated the trajectory of the (unitary) synaptic conductance in one cycle by a simple triangle (Fig. 5A), which could be defined by three parameters: duration ($T_{\text{act}}$), peak time ($t_{\text{peak}}$) and amplitude ($g_{\text{max}}$) (Fig. 5B). This simplified triangular synaptic conductance waveform could then be repeated with any period ($P$) to mimic the realistic synaptic input to the LP neuron. For a given synaptic duration $T_{\text{act}}$, the peak phase of the synapse can be defined as $\Delta_{\text{peak}} = t_{\text{peak}} / T_{\text{act}}$. The parameter $\Delta_{\text{peak}}$ is known to vary as a function of $P$ (Tseng et al., 2014) and, in a previous study, we found that $\Delta_{\text{peak}}$ may influence the activity of the postsynaptic neuron, independent of $P$ and $g_{\text{max}}$ (Mamiya and Nadim, 2004). We therefore systematically explored the influence of three parameters of the synaptic input ($P$, $g_{\text{max}}$ and $\Delta_{\text{peak}}$) on $\phi_{\text{LP}}$.

As with the realistic synaptic waveforms (Fig. 3), we used the dynamic clamp technique to apply the triangular conductance waveform periodically to the LP neuron in the presence of the synaptic blocker picrotoxin. Across different runs within the same experiment, the parameters $P$, $g_{\text{max}}$ and $\Delta_{\text{peak}}$ were changed on a grid (see Methods). In addition, all combinations of these three parameter values were run in two conditions in the same experiment, 1: with constant duration, i.e., constant $T_{\text{act}}$ across different $P$ values (C-Dur of 300 ms), and 2: with constant duty cycle, i.e., $T_{\text{act}}$ changing
proportionally to $P$ (C-DC of 0.3; Fig. 5C). Using these protocols, we measured the effects of synaptic parameters on $\phi_{LP}$ (Fig. 5D).

The LP neuron produced burst responses that followed the synaptic input in a 1:1 manner across all values of $P$ that were used (Fig. 6A1). When $g_{\text{max}}$ and $\Delta_{\text{peak}}$ were kept constant, $\phi_{LP}$ decreased as a function of $P$ (Fig. 6A2). This decrease was always larger for the C-Dur case than the C-DC case. For both C-DC and C-Dur, this trend was seen across all values of $\Delta_{\text{peak}}$ and $g_{\text{max}}$ (Fig 6A3). The effect of $P$ on $\phi_{LP}$ was highly significant for both C-DC (Three-Way ANOVA, $p<0.001$, $F=100.677$) and C-Dur (Three-Way ANOVA, $p<0.001$, $F=466.424$), indicating that the period and duration of the inhibitory input to the LP neuron had a significant effect on its phase.

Changing $g_{\text{max}}$ produced a large effect on the level of hyperpolarization in the LP neuron, but this usually translated to only a small or modest effect on the time to the first spike following inhibition (Fig. 6B1). Overall, increasing $g_{\text{max}}$ at constant values of $P$ and $\Delta_{\text{peak}}$ produced a significant but only small to moderate increase in $\phi_{LP}$ (Three-Way ANOVA, $p<0.001$, $F=10.798$). Although increasing $g_{\text{max}}$ produced the same qualitative effect for both the C-DC and C-Dur (e.g., Fig. 6B2), $\phi_{LP}$ in the C-DC case was restricted to a smaller range (Fig. 6B3 top vs. bottom panels). Overall, this increase was robust for most values of $P$ and $\Delta_{\text{peak}}$ (Fig. 6B3).

Increasing $\Delta_{\text{peak}}$ for a constant value of $P$ and $g_{\text{max}}$ (Fig. 6C1), produced a small but significant increase in $\phi_{LP}$ (Three-Way ANOVA, $p<0.001$, $F=17.172$). This effect was robust for most values of $P$ and $g_{\text{max}}$, for both C-DC and C-Dur (Fig. 6C2 and 6C3).

These results showed that all three parameters that define the shape of the IPSC influence $\phi_{LP}$. Clearly, the strongest effect is the decrease in $\phi_{LP}$ as a function of $P$. However, $\phi_{LP}$ modestly increases as a function of the other two parameters, $g_{\text{max}}$ and $\Delta_{\text{peak}}$. This raised the question how $g_{\text{max}}$ and $\Delta_{\text{peak}}$ would have to change in coordination as a function of $P$ to counteract the effect of $P$ on $\phi_{LP}$ and achieve phase constancy.

**Coordinated changes of $g_{\text{max}}$ and $\Delta_{\text{peak}}$ produce the largest effect on phase**

To explore how $g_{\text{max}}$ and $\Delta_{\text{peak}}$ might interact to influence $\phi_{LP}$, we examined the sensitivity of $\phi_{LP}$ to these two parameters, individually and in combination, for all values of $P$ in our data (see Methods). Sensitivity of $\phi_{LP}$ to these two parameters varied across $P$ values, with larger sensitivity at lower values of $P$ (Two-Way RM-ANOVA, $p<0.001$,
F=16.054; data included in Figure 7-source data). For simplicity, we averaged the sensitivity values across different $P$ values to obtain an overall measure of the influence of $g_{\text{max}}$ and $\Delta_{\text{peak}}$. These results showed that, for the C-DC case, $\phi_{\text{LP}}$ had a positive sensitivity to $g_{\text{max}}$ and a smaller positive sensitivity to $\Delta_{\text{peak}}$ (Fig. 7A). The sensitivity was largest if the two parameters were varied together ($g_{\text{max}} + \Delta_{\text{peak}}$) and smallest if they were varied in opposite directions ($g_{\text{max}} - \Delta_{\text{peak}}$; Two-Way RM-ANOVA, $p<0.001$, $F=3.330$). Similarly, these sensitivity values were also significantly different for the C-Dur case (Fig. 7B; Two-Way RM-ANOVA, $p<0.001$, $F=2.892$), with largest sensitivity for $g_{\text{max}} + \Delta_{\text{peak}}$ and smallest for $g_{\text{max}} - \Delta_{\text{peak}}$.

**Level sets of $\phi_{\text{LP}}$ in the $P$-$g_{\text{max}}$-$\Delta_{\text{peak}}$ space for C-DC and C-Dur cases**

To search for phase constancy across different $P$ values in our dataset, we expressed $\phi_{\text{LP}}$ as a function of the three IPSC parameters, $P$, $g_{\text{max}}$ and $\Delta_{\text{peak}}$:

$$\phi_{\text{LP}} = \Phi(P, g_{\text{max}}, \Delta_{\text{peak}}).$$

Figure 8 shows heat map plots of the function $\Phi$, plotted for the range of values of $P$ and $\Delta_{\text{peak}}$ and four values of $g_{\text{max}}$. In these plots, phase constancy can be seen as the set of values in each graph that are isochromatic, indicating the level sets of the function $\Phi$. These level sets are mathematically defined as hypersurfaces on which the function has a constant value: $\Phi(P, g_{\text{max}}, \Delta_{\text{peak}}) = \phi_c$. For the C-DC case, in each $g_{\text{max}}$ section of the plot, the level sets (e.g. $\phi_c = 0.34$ denoted in white) spanned a moderate range of $P$ values as $\Delta_{\text{peak}}$ increased (Fig. 8A1). The span of $P$ values across all four panels indicates the range of cycle periods for which phase constancy could be achieved by varying $g_{\text{max}}$ and $\Delta_{\text{peak}}$. This range of $P$ values (spanned by the white curves) was considerably smaller for the C-Dur case (Fig. 8A2).

For any constant phase value $\phi_c$, these level sets can be expressed as

$$P = P_{\phi_c}(g_{\text{max}}, \Delta_{\text{peak}}),$$

which describes a surface in the 3D space, yielding the $P$ value for which phase can be maintained at $\phi_c$, for the given values of $g_{\text{max}}$ and $\Delta_{\text{peak}}$. The level set indicated by the white curves in panel A for the C-DC case is plotted as a heat map in Fig. 8B1 and can be compared with the same plot for the C-Dur case in Fig. 8B2. The range of colors in each plot (marked next to each panel) indicates the range of $P$ values for which phase can be kept at $\phi_c = 0.34$. To reveal how this range depends on the desired phase, we measured this range for all values of $\phi_c$ between 0.2 and 0.8 (Figs. 8C1 and 8C2). We
found that the LP neuron could not achieve phases below 0.3 in the C-DC case (Fig. 8C1), which is simply because the neuron never fired during the inhibitory synaptic current (which had a duty cycle of 0.3). Furthermore, the range of $P$ values for which the LP phase could be maintained by varying $g_{\text{max}}$ and $\Delta_{\text{peak}}$ was much larger for C-DC inputs compared to C-Dur Inputs, for all $\varphi_c$ values between 0.31 and 0.54.

A model of synaptic dynamics could predict activity onset phase of the LP neuron

To gain a better understanding of our experimental results, we derived a mathematical description of the phase of a follower neuron such as LP, based on the following assumptions: 1, that the firing time of this neuron was completely determined by its synaptic input, 2, that in each cycle the synaptic conductance $g_{\text{syn}}$ increased to a maximum value $g_{\text{max}}$ for a time interval $T_{\text{act}}$ (the active duration of the synapse) and decayed to 0 otherwise, and 3, that the follower neuron remained inactive when $g_{\text{syn}}$ was above some threshold $g^*$. The derivation of this model is described in the Methods.

This simple model provided a mathematical description of $\varphi_{LP}$ as a function of $P$, $g_{\text{max}}$ and $\Delta_{\text{peak}}$, for the C-Dur and C-DC cases. In the C-Dur case (Equation (7)), as $P$ increased, $\varphi_{LP}$ decayed and approached 0 like $1/P$. In contrast, in the C-DC case (Equation (8)), $\varphi_{LP}$ approached its lower limit $\Delta_{\text{peak}} \cdot \text{DC}$, as $P$ increased, and thus behaved very differently than in the C-Dur case.

We used these equations to describe $g_{\text{max}}$ as a function of $P$ (for any given $\Delta_{\text{peak}}$) so that LP maintained a constant phase $\varphi_c$, (Equation (10) for the C-DC case). Alternatively, $\Delta_{\text{peak}}$ could be given as a function of $P$ (for any given $g_{\text{max}}$, Equation (11) for the C-DC case). We used these derivations to compare how phase constancy depends on $g_{\text{max}}$ or $\Delta_{\text{peak}}$ in the C-DC case. A comparison of these two cases can be seen in Fig. 9A, where either $g_{\text{max}}$ (green) or $\Delta_{\text{peak}}$ (blue) is varied to keep $\varphi_{LP}$ constant at $\varphi_c=0.34$ across different $P$ values. (The red curve is the depressing case, described below.) As the figure shows, phase constancy can be achieved by varying either parameter, but each parameter produces a different range of $P$ across which phase is maintained.

These equations and their corresponding counterparts for the C-Dur case can be used to calculate the range of $P$ values over which changing $\Delta_{\text{peak}}$ (from 0 to 1) can maintain a constant phase $\varphi_c$. If $\Delta P$ denotes the range of $P$ values for which phase can
be constant, it is straightforward to show that $\Delta P_{DC} > \Delta P_{Dur}$ (compare black and blue curves in Fig. 9B and 9C; see Methods for derivation).

Two additional points are notable in Fig. 9C. First, the lower bound on $\phi_{LP}$ for which phase constancy can occur is smaller in the C-Dur than the C-DC case. This is because we have assumed that in the C-DC case the LP neuron cannot fire during inhibition (i.e., until after $\Delta_{peak} DC$). Second, for $\phi_c$ larger than $\sim 0.5$, $\Delta P$ is larger for the C-Dur case. This occurs because Equation (12) can no longer be satisfied when $\phi_c$ is large. That is, with constant duty cycle, it is not possible to produce an arbitrarily large follower neuron phase, but with constant duration, any large phase is attainable if the cycle period is not much larger than the synaptic duration. These findings are consistent with our experimental results described above (see Fig. 8).

The pacemaker synaptic input to the LP neuron shows short-term synaptic depression (Rabbah and Nadim, 2007). In a previous modeling study, we explored how the phase of a follower neuron was affected when the inhibitory synapse from an oscillatory neuron to this follower had short-term synaptic depression (Manor et al., 2003). In that study the role of the parameter $\Delta_{peak}$ was not considered. We now consider how the presence of short-term synaptic depression influences phase constancy by changing both $g_{max}$ and $\Delta_{peak}$. As stated in the Methods (Equation (16)), the effect of synaptic depression on synaptic strength can be obtained as $g_{max} = \tilde{g}_{max} \cdot s_{max}(P)$, where $s_{max}$ is an increasing function whose value approaches 1 as $P$ increases. This indicates that the synapse becomes stronger due to more recovery from depression at longer cycle periods. When synaptic depression dictates how $g_{max}$ varies with $P$ and $\Delta_{peak}$ also varies with $P$ and $g_{max}$ (Equation (11)), the simultaneous changes in $g_{max}$ and $\Delta_{peak}$ (red) greatly increase the range of $P$ values over which $\phi_{LP}$ is constant (Fig. 9A).

Note that the C-DC case with short-term depression spans a larger range of $P$ values than the non-depressing case (Fig. 9B). Similarly, the range of $P$ values for which phase can be maintained is larger than the non-depressing case across $\phi_{LP}$ values, except where $\phi_{LP}$ is so large that the depressing synapse operates outside its dynamic range (Fig. 9C). These results are consistent with our experimental results, indicating that although phase constancy can be achieved when either $g_{max}$ or $\Delta_{peak}$ increases with $P$, a concomitant increase of both - which could occur for example with
Discussion

The importance of phase in oscillatory networks

A common feature of oscillatory networks is that the activities of different neuron types are restricted to specific phases of the oscillation cycle. For example, different hippocampal and cortical neurons are active in at least three distinct phases of the gamma rhythm (Hajos et al., 2004; Hasenstaub et al., 2005), and distinct hippocampal neuron types fire at different phases of the theta rhythm and sharp wave-associated ripple episodes (Somogyi and Klausberger, 2005).

Experimental studies quantify the latency of neural activity with respect to a reference time in the cycle, but in most cases, these latencies are normalized and reported as phase. Distinct neuron types can maintain a coherent activity phase, despite wide variations in the network frequency (30-100 Hz for gamma rhythms, 4-7 Hz for theta rhythms, and 120-200 Hz for sharp wave-associated ripple episodes). Phase-specific activity of different neuron types is proposed to be important in rhythm generation (Wang, 2010), and indicates the necessity of precise timing for producing proper circuit output and behavior (Kopell et al., 2011). For example, phase locking of spike patterns to oscillations is important for auditory processing, single cell and network computations and Hebbian learning rules (Kayser et al., 2009; McLelland and Paulsen, 2009; Panzeri et al., 2010). For brain oscillations, phase relationships may provide clues about the underlying circuit connectivity and dynamics, but a behavioral correlate of varying frequencies is not obvious. In contrast, the activity phase of distinct neuron types in rhythmic motor circuits is a tangible readout of the timing of motor neurons and muscle contractions, thus defining phases of movement (Grillner and El Manira, 2015; Kiehn, 2016; Le Gal et al., 2017; Bidaye et al., 2018). Because meaningful behavior depends crucially on proper activity phases, whether neurons maintain their activity phase in face of changes in frequency simply translates to whether the movement pattern changes as it speeds up or slows down.

Determinants of phase

In oscillatory networks, the activity phases of different neuron types depend to different degrees on the precise timing and strength of their synaptic inputs (Oren et
Our results from noise current injections showed that the timing of the LP neuron is strongly dependent on the timing of inputs it receives. Dynamic clamp injection of realistic or triangular conductance waveforms with different periods ($P$) indicated that $\phi_{LP}$ was largely determined by the duration of the synaptic input. $\phi_{LP}$ changed substantially with $P$ when inputs had constant duration, but much less when inputs had a constant duty cycle, i.e., when duration scaled with $P$. However, our experiments also showed that inputs of constant duty cycles alone are insufficient for phase constancy. $\phi_{LP}$ decreased with $P$ even with a constant duty cycle of inputs, but increased with either synaptic strength ($g_{max}$) or peak phase of the synaptic input ($\Delta_{peak}$). The increase in $\phi_{LP}$ had similar sensitivity to $g_{max}$ and $\Delta_{peak}$, and therefore a larger sensitivity to a simultaneous increase in both. Consequently, it was possible to keep $\phi_{LP}$ constant over a wide range of cycle periods by increasing both parameters with $P$.

The fact that an increase in $g_{max}$ with $P$ promotes phase constancy is biologically relevant, as short-term depression in pyloric synapses means that synaptic strength indeed increases with $P$ (Manor et al., 1997). Previous modeling studies show that short-term synaptic depression of inhibitory synapses promotes phase constancy (Nadim et al., 2003; Bose et al., 2004), largely because of longer recovery times from depression at larger values of $P$.

The finding that an increase of $\Delta_{peak}$ with $P$ promotes phase maintenance is somewhat surprising, as we have previously shown that $\Delta_{peak}$ in LP actually decreases with $P$ (Manor et al., 1997; Tseng et al., 2014). On the face of it, this suggests that an increase in $\Delta_{peak}$ is not a strategy employed in the intact circuit. However, the caveat is that such results may critically depend on the cause of the change in $P$, either technically and biologically. While in our current study we varied $\Delta_{peak}$ with direct conductance injection into LP, previous results were obtained by changing the waveform and period of the presynaptic pacemaker neurons. When $P$ is changed in an individual preparation by injecting current into or voltage-clamping the pacemakers, phase of follower neurons is not particularly well maintained. An example of this is shown in Fig. 1, where $\phi_{LP}$ values fall between constant phase and constant duration and, additionally, all pyloric neurons show behavior that falls between constant phase and constant latencies (Hooper, 1997b, a). This may reflect that neurons are not keeping phase particularly well when the only cause of changing $P$ is the presynaptic
input. This is supported by the observation that even during normal ongoing pyloric activity, phases change with cycle-to-cycle variability of $P$ in individual preparations (Bucher et al., 2005). However, it does not preclude the possibility that $\Delta_{\text{peak}}$ plays an important role in stable phase relationships when $P$ differs because of temperature, neuromodulatory conditions, or inter-individual variability (discussed below).

It is noteworthy that a change in the synaptic strength or peak phase with $P$ is not peculiar to graded synapses. The fact that short-term synaptic plasticity can act as a frequency-dependent gain control mechanism is well known for many spike-mediated synaptic connections. In bursting neurons, the presence of a combination of short-term depression and facilitation in the same spike-mediated synaptic interaction could also result in changes in the peak phase of the summated synaptic current as a function of burst frequency and duration, and the intra-burst spike rate (Markram et al., 1998).

The mathematical model in the current study provides mechanistic explanations for several of our experimental findings. First, it can be used to produce a quantitative measure of phase, given the values of $g_{\text{max}}$, $\Delta_{\text{peak}}$ and $P$. Thus, these equations can be used to compare the C-DC and C-Dur cases, which match our experimental results. They show that, for most phase values, the C-DC case provides a larger range of cycle periods at which phase constancy can occur. Second, these equations provide the activity phase no matter how the pacemaker synaptic input duration changes with cycle period. For instance, our experiments were conducted by changing synaptic input through sampling individual values of the parameter pairs $g_{\text{max}}$ and $\Delta_{\text{peak}}$, and then calculating the ensuing phase. We then used fitting to find level sets of constant phase (Fig. 8). In contrast, when we combined our mathematical derivation here with previous results on the role of short-term synaptic depression (Bose et al., 2004), we could demonstrate how a neuronal circuit can naturally follow a level set of phase (Equations (7), (8), (15) and (16)). Moreover, we showed that the combined increase in $g_{\text{max}}$ and $\Delta_{\text{peak}}$ with $P$ produces a larger range of periods for phase constancy than increasing either parameter alone. In short, this mathematical formulation produces a simple quantitative distillation of our experimental results.

In this study, we did not explore the role of the intrinsic properties of the LP neuron on its phase. In separate experiments, we simultaneously measured post-inhibitory rebound properties in synaptically isolated LP neurons and the levels of voltage-gated ionic currents (the transient potassium current $I_A$ and the
hyperpolarization-activated inward current \(I_h\) that influence rebound spiking. These data were not included in this study for brevity and because they showed that the timing of post-inhibitory spiking was relatively stable across preparations. Therefore, we would expect the contribution of intrinsic properties in controlling the timing of the LP neuron burst onset to be relatively small. However, this result does not generalize to all follower neurons. For example, the follower ventral dilator (VD) and PY neurons have a much higher levels of \(I_h\), which in turn has a larger effect on the timing of post-inhibitory spiking. In a set of computational studies, we addressed the role of \(I_h\) in determining the burst phase in response to periodic inputs (Zhang et al., 2008, 2009) and in conjunction with short-term depression in the synaptic input (Bose et al., 2004). An experimental clarification of the relative contribution of intrinsic properties vs. synaptic input could be done with controlled dynamic clamp synaptic input, such as those used in the current study, injected in PY or VD neurons. Such a data set would fittingly complement the results of the current study to elucidate more general rules in determining the activity phase of neurons in an oscillatory network.

Phase relationships in changing temperatures

An interesting case is provided by the observation that phases are remarkably constant when pyloric rhythm frequency is changed with temperature. Tang et al. (2012) report a 4-fold decrease in \(P\) of the pyloric rhythm between 7 and 23° C. In this study, none of the pyloric phases changed significantly, and it is worth noting that under conditions of changing temperatures, the relationships between \(P\), \(g_{\text{max}}\), and \(\Delta_{\text{peak}}\) appeared to be fundamentally different from when \(P\) is changed at a constant temperature. Presynaptic voltage trajectories scaled with changing \(P\), and \(\Delta_{\text{peak}}\) of postsynaptic currents was independent of \(P\), in contrast to the decrease described at constant temperature (Manor et al., 1997; Tseng et al., 2014). Amplitudes of synaptic potentials did not change with temperature, despite an increase in synaptic current amplitudes with increasing temperature (and associated decrease in \(P\)). This is in contrast to the positive relationship between \(g_{\text{max}}\) and \(P\) that results from synaptic depression at a constant temperature (Manor et al., 1997). Therefore, it appears that the likely substantial effects of temperature on synaptic dynamics and ion channel gating are subject to a set of compensatory adaptations different from when \(P\) is changed at constant temperature.

Variability and slow compensatory regulation of phase
Phase maintenance in the face of changing $P$ in an individual animal requires the appropriate short-term dynamics of synaptic and intrinsic neuronal properties. The fact that characteristic (and therefore similar) phase relationships can also be observed under the same experimental conditions across individual preparations is a different conundrum, particularly when $P$ can vary substantially, as is true for brain oscillations (Hajos et al., 2004; Hasenstaub et al., 2005; Somogyi and Klausberger, 2005). Phases show different degrees of variability across individuals in a variety of systems, e.g., leech heartbeat (Wenning et al., 2018), larval crawling in *Drosophila* (Pulver et al., 2015), and fictive swimming in zebrafish (Masino and Fetch, 2005), but in all of these cases phases are not correlated with $P$. In the pyloric rhythm, phases are also variable to a degree across individuals, but not correlated with the mean $P$, which varies >2-fold (Bucher et al., 2005; Goaillard et al., 2009). This phase constancy occurs despite considerable inter-individual variability in ionic currents, and is considered the ultimate target of slow compensatory regulation, i.e., homeostatic plasticity (Marder and Goaillard, 2006; Ma and LaMotte, 2007; Marder et al., 2014). Slow compensation can also be observed directly when rhythmic activity is disrupted by decentralization, and subsequently recovers to similar phase relationships over the course of many hours (Luther et al., 2003). It is interesting to speculate if our findings about how synaptic parameters must change to keep phase constant would hold across individuals with different mean $P$. The prediction would be coordinated positive correlations of both $g_{\text{max}}$ and $\Delta_{\text{peak}}$ with $P$.

Synaptic inputs to the LP neuron show considerable variability across preparations (e.g., Fig. 3B), which mirrors the variability seen in the levels of voltage-gated ionic currents in pyloric neurons (Schulz et al., 2006). We did not address the role and extent of variability in this study, because a proper analysis of variability required us to first establish the mechanisms that give rise to a consistent output, in this case phase constancy. Based on our findings regarding the influence of synaptic parameters on phase, a natural next step is to explore whether the variability of different parameters defining the synaptic input influences variability of phase or, alternatively, whether variability in some synaptic parameters may be irrelevant to phase or restrained by the postsynaptic neuron.

*Phase relationships under different neuromodulatory conditions.*
The flipside of the question how neurons maintain phase is the question how their phase can be changed. In motor systems in particular, changes in phase relationships are functionally important to produce qualitatively different versions of circuit output, for example to produce different gaits in locomotion (Vidal-Gadea et al., 2011; Grillner and El Manira, 2015; Kiehn, 2016). The activity of neural circuits is flexible, and much of this flexibility is provided by modulatory transmitters and hormones which alter synaptic and intrinsic neuronal properties (Brezina, 2010; Harris-Warrick, 2011; Jordan and Slawinska, 2011; Bargmann, 2012; Marder, 2012; Bucher and Marder, 2013; Nadim and Bucher, 2014). The pyloric circuit is sensitive to a plethora of small molecule transmitters and neuropeptides which affect cycle frequency and phase relationships (Marder and Bucher, 2007; Stein, 2009; Daur et al., 2016). Indeed, extensive research has indicated the role of amine modulation of synaptic strength and neuronal firing phase in the pyloric circuit, and how amine modulation of synaptic and intrinsic firing properties changes firing phases (Johnson et al., 2003; Gruhn et al., 2005; Johnson et al., 2005; Peck et al., 2006; Harris-Warrick and Johnson, 2010; Harris-Warrick, 2011; Kvarta et al., 2012). With respect to our findings, any given neuromodulator could act presynaptically to alter $P$, duration, or duty cycle on the one hand, and $g_{\text{max}}$ and $\Delta_{\text{peak}}$ on the other. In addition, the neuromodulator could affect the postsynaptic neuron’s properties and alter its sensitivity to any of these parameters. Therefore, our findings could not just further our understanding of how phase can be maintained across different rhythm frequencies, but also provide a framework for testing if and how changes in synaptic dynamics may contribute to altering phase relationships under different neuromodulatory conditions.

**Materials and Methods**

Adult male crabs (*Cancer borealis*) were acquired from local distributors and maintained in aquaria filled with chilled (10-13°C) artificial sea water until use. Crabs were prepared for dissection by placing them on ice for 30 minutes. The dissection was performed using standard protocols as previously described (Tohidi and Nadim, 2009; Tseng and Nadim, 2010). The STNS, including the four ganglia (esophageal ganglion, two commissural ganglia, and the STG) and their connecting nerves, and the motor nerves arising from the STG, were dissected from the stomach and pinned into a Sylgard (Dow-Corning) lined Petri dish filled with chilled saline. The STG was desheathed, exposing the somata of the neurons for intracellular impalement.
Preparations were superfused with chilled (10-13°C) physiological Cancer saline containing: 11 mM KCl, 440 mM NaCl, 13 mM CaCl$_2$ $\cdot$ 2H$_2$O, 26 mM MgCl$_2$ $\cdot$ 6H$_2$O, 11.2 mM Trizma base, 5.1 mM maleic acid with a pH of 7.4.

Extracellular recordings were obtained from identified motor nerves using stainless steel electrodes, amplified using a differential AC amplifier (A-M Systems, model 1700). One lead was placed inside a petroleum jelly well created to electrically isolate a small section of the nerve, the other right outside of it. For intracellular recordings, glass microelectrodes were prepared using the Flaming-Brown micropipette puller (P97; Sutter Instruments) and filled with 0.6 M K$_2$SO$_4$ and 20 mM KCl. Microelectrodes used for membrane potential recordings had resistances of 25-30MΩ; those used for current injections had resistances of 15-22 MΩ. Intracellular recordings were performed using Axoclamp 2B and 900A amplifiers (Molecular Devices). Data were acquired using pClamp 10 software (Molecular Devices) and the Netsuite software (Gotham Scientific), sampled at 4-5 kHz and saved on a PC using a Digidata 1332A (Molecular Devices) or a PCI-6070-E data acquisition board (National Instruments).

Individual pyloric neurons were impaled and identified via their membrane potential waveforms, correspondence of spike patterns with extracellular nerve recordings, and interactions with other neurons within the network (Weimann et al., 1991).

**Constructing realistic graded IPSC waveforms**

Inhibitory postsynaptic currents (IPSCs) were recorded from LP neurons during the ongoing rhythm using two-electrode voltage clamp and holding the LP neuron at -50mV, far from the IPSC reversal potential of ~ -80 mV (Fig. 3A). We refer to the total current measured in the voltage clamped LP neuron during the activity of the PD and PY neurons as a synaptic current for the following reasons: 1, the synaptically isolated LP neuron produces tonic spiking activity (see, e.g., Fig. 2B), and 2, holding the LP neuron at different voltages (e.g. -60 or -110 mV) produces a similarly shaped current, but with a different amplitude or reversed sign (at -110 mV).

When the LP soma is voltage clamped at -50 mV, the axon (which is electrotonically distant from the soma) produced action potentials following the synaptic inhibition from the PY neuron and the pacemaker neurons. The onset of the LP neuron
action potentials (recorded in the current trace) was used to calculate the mean IPSC for each experiment averaging the IPSCs over 10-20 cycles. The IPSC waveforms were then extracted by normalizing both the amplitude and the duration of the mean IPSC.

**Driving the LP neuron with noise current**

In these experiments, the preparation was superfused in *Cancer* saline plus 10⁻⁵ M picrotoxin (PTX; Sigma Aldrich) for 30 minutes to block the synaptic currents to the LP neuron. The removal of synaptic inhibition onto LP neurons changed the activity of these neurons from bursting to tonic firing. Then, noise current, generated by the Ornstein-Uhlenbeck (O-U) process (Lindner), was injected into the LP neurons for 60 minutes using the Scope software (available at [http://stg.rutgers.edu/Resources.html](http://stg.rutgers.edu/Resources.html), developed in the Nadim laboratory). The baseline of the noise current was adjusted by adding DC current so that it can provide enough inhibition to produce silent periods alternating with bursts of action potentials. The O-U process was defined as

\[ dX_t = -\frac{1}{\tau} X_t dt + \sigma dW_t. \]

The parameters used for noise injection were \( \tau = 10 \) to 20 ms, \( \sigma = 200 \) pA and a DC current of -200 to -100 pA. In these experiments we defined bursts as groups of at least two action potentials with inter-spike intervals < 300 ms, following a gap of at least 300 ms.

**Driving the LP neuron with realistic or triangular IPSC waveforms in dynamic clamp**

The dynamic clamp current was injected using the Netclamp software (Netsuite, Gotham Scientific). We pharmacologically blocked synaptic inputs from the pacemaker AB and follower PY neurons to the LP neuron by superfusing the preparation in *Cancer* saline plus 10⁻⁵ M picrotoxin (PTX; Sigma Aldrich) for 30 minutes. This treatment does not block the cholinergic synaptic input from the PD neurons for which no clean pharmacological blocker is known. Although the PD neuron input has some influence on the LP neuron activity, this input only constitutes <20% of the total pacemaker synapse and cannot drive oscillations in the follower LP neuron.
The LP neuron was driven in PTX with an artificial synaptic current in dynamic clamp. The synaptic current was given as
\[
I_{syn} = g_{syn}(V_{LP} - E_{syn})
\]
where the synaptic conductance \(g_{syn}\) was a pre-determined waveform, repeated periodically with period \(P\), and \(E_{syn}\) was the synaptic reversal potential set to \(-80\) mV (Zhao et al., 2011).

Two sets of dynamic clamp experiments were performed on different animals. In one set of experiments, \(g_{syn}\) was set to be a triangular waveform. We measured the effects of four different parameters in these triangle conductance injections (Fig. 1): peak phase (\(\Delta_{peak}\)), duration (\(T_{act}\)), period (\(P = \) time between onsets of dynamic clamp synaptic injections), and maximal conductance (\(g_{max}\), the peak value of \(g_{syn}\)). This allowed us to explore which combinations of the different parameters influences the LP phase. Five values for \(P\) were used: 500, 750, 1000, 1500, and 2000 ms, which cover the typical range of pyloric cycle periods. Three values of \(g_{max}\) were used: 0.1, 0.2 and 0.4 \(\mu\)S, consistent with previous measurements of synaptic conductance (Zhao et al., 2011; Tseng et al., 2014). The value of \(\Delta_{peak}\) was varied to be 0, 0.25, 0.5, 0.75 or 1. In the same experiment, all runs were done in two conditions: with \(T_{act}\) constant across different \(P\) values (C-Dur case with \(T_{act} = 300\) ms) or with \(T_{act}\) changing proportionally to \(P\) (C-DC case with duty cycle \(DC = T_{act}/P = 0.3\)).

In the other set of experiments, \(g_{syn}\) was a realistic IPSC waveform, based on a pre-recorded IPSC in the LP neuron. In these experiments, \(P\) was varied to be 500, 750, 1000, 1250, 1500, or 2000 ms by scaling the realistic waveform in the time direction. In these experiments, \(g_{max}\) was set to be 0.1, 0.2, 0.4, 0.6, or 0.8 \(\mu\)S. The LP neuron burst onset delay (\(\Delta t\)) was measured relative to the onset of the pacemaker component of the synaptic input (identified by the kink in the synaptic conductance waveform) in each cycle. The burst phase was calculated as \(\phi_{LP} = \Delta t / P\). Phase constancy means that \(\Delta t\) changed proportionally to \(P\). To measure the LP neuron phase with respect to the end of the pacemaker input, this reference used was the point on the synaptic conductance waveform marked by drawing a horizontal line from the kink that identified the onset of the pacemaker input.
Determining relationship between cycle period (P), synaptic strength (g_{max}) and LP phase (\varphi_{LP}) using the realistic IPSC waveform

We determined how well the mathematical model derived for constant input duty cycles (see Equation (8) below), matched the experimental data obtained with realistic IPSC waveforms. To this end, we fit the model to \varphi_{LP} values measured for all values of g_{max} and P, using the standard fitting routine 'fit' in MATLAB (Mathworks).

**Sensitivity of \varphi_{LP} to g_{max} and \Delta_{peak} across all P values**

To explore how g_{max} and \Delta_{peak} may interact to influence \varphi_{LP}, we examined the sensitivity of \varphi_{LP} to these two parameters, individually and in combination, for all values of P in our data. For each P, we computed the mean value of \varphi_{LP} across all experiments, and all values of g_{max} (0.1, 0.2, 0.3 and 0.4 \mu S) and \Delta_{peak} (0, 0.25, 0.5, 0.75 or 1). (The \varphi_{LP} value for g_{max} = 0.3 \mu S was obtained in this case by linearly interpolating the values for 0.2 and 0.4 \mu S.) This produced a 4 by 5 matrix of all values.

For each data point in the matrix, we moved along eight directions (+g_{max}, +\Delta_{peak}, -g_{max}, -\Delta_{peak}, +g_{max} & +\Delta_{peak}, -g_{max} & -\Delta_{peak}, +g_{max} & -\Delta_{peak}, +\Delta_{peak} & -\Delta_{peak}). Here “+” denotes increasing and “−” denotes decreasing. We then calculated the change in \varphi_{LP} per unit g_{max} (normalized by 0.4 \mu S, \Delta_{peak}, or both. For example, the sensitivity of \varphi_{LP} when \Delta_{peak} was changed from 0.25 to 0.5 was measured as

\[
\frac{\varphi_{LP} \text{ (at } \Delta_{peak} = 0.5) - \varphi_{LP} \text{ (at } \Delta_{peak} = 0.5)}{0.5 - 0.25}
\]

Similar, the sensitivity of \varphi_{LP} when g_{max} was changed from 0.2 to 0.4 was measured as

\[
\frac{\varphi_{LP} \text{ (at } g_{max} = 0.4) - \varphi_{LP} \text{ (at } g_{max} = 0.2)}{(0.4 - 0.2) / 0.4}
\]

These data are provided in Figure 7-source data. As the next step, we averaged the sensitivity along each aligned direction: [+g_{max} and -g_{max}]; [+\Delta_{peak} and -\Delta_{peak}]; [+g_{max} & +\Delta_{peak} and -\Delta_{peak}]; [+g_{max} & -\Delta_{peak} and +g_{max} & -\Delta_{peak}]. This produced the four cardinal directions, shown in Fig. 7. Finally, we averaged the sensitivity across all P values.

**A model of synaptic dynamics**
In the derivation of the model, the firing time of the LP neuron was assumed to be completely determined by its synaptic input. This synaptic conductance \( g_{\text{syn}} \) was assumed to rise and fall with distinct time constants. The following holds over one cycle period and therefore time is reset with period \( P \) \((t \mod P)\):

\[
\frac{dg_{\text{syn}}}{dt} = \begin{cases} (g_{\text{max}} - g_{\text{syn}}) \tau_r, & t \mod P < t_{\text{peak}} \\ -g_{\text{syn}} / \tau_s, & t \mod P \geq t_{\text{peak}} \end{cases}
\]  

(1)

where the time \( t_{\text{peak}} \), corresponding to \( \Delta_{\text{peak}} \), is \( t_{\text{peak}} = \Delta_{\text{peak}} T_{\text{act}} \). We assumed that LP neuron remained inactive when \( g_{\text{syn}} \) was above a fixed threshold \( (g^*) \) less than \( g_{\text{max}} \).

Because the synaptic input is periodic with period \( P \), we solved for the minimum and maximum values of \( g_{\text{syn}} \) in each cycle. The minimum \( (g_{\text{lo}}) \) occurred just before the onset \((t = 0)\) of AB/PD activity, whereas the maximum occurred at the peak synaptic phase \( \Delta_{\text{peak}} \) for the C-Dur case. In the C-DC case, \( T_{\text{act}} = DC \cdot P \), where DC is the duty cycle (fixed at 0.3 in our experiments).

To calculate \( g^* \), we set the value \( t = 0 \) so that \( g_{\text{syn}}(0) = g_{\text{lo}} \) (and, by periodicity, \( g_{\text{syn}}(P) = g_{\text{lo}} \)), and solved the first part of Equation (1) where \( g_{\text{syn}} \) increases until \( t = t_{\text{peak}} \).

This yielded

\[
g_{\text{peak}} = g_{\text{syn}}(t_{\text{peak}}) = g_{\text{max}} + (g_{\text{lo}} - g_{\text{max}}) e^{-t_{\text{peak}} / \tau_s}
\]

(2)

We then used the second part of Equation (1) to track the decay of \( g_{\text{syn}} \) for \( t_{\text{peak}} < t < P \):

\[
g_{\text{syn}}(t) = g_{\text{peak}} e^{-(t-t_{\text{peak}}) / \tau_s}
\]

(3)

Using Equation (3), we calculated the time \( \Delta t \) at which the synaptic conductance \( g_{\text{syn}}(\Delta t) = g^* \) as follows:

\[
g^* = g_{\text{peak}} e^{-\left(\Delta t - t_{\text{peak}}\right) / \tau_s}
\]

(4)

Solving Equation (4) for \( \Delta t \) yielded

\[
\Delta t = \tau_{\text{s}} \ln \frac{g(t_{\text{peak}})}{g^*} + t_{\text{peak}}.
\]

Dividing this equation by \( P \) yielded \( \varphi_{LP} \):
\[ \phi_{LP} = F(P, g_{\text{max}}, \Delta_{\text{peak}}) = \frac{\tau_s}{P} \ln \frac{g_{\text{peak}}}{g^*} + \frac{t_{\text{peak}}}{P}, \]  

(5)

where \( g_{\text{peak}} \) is given by Equation (2). This expression provides a description of the dependence of \( \phi_{LP} \) as a function of \( P, g_{\text{max}} \) and \( \Delta_{\text{peak}} \). To explore the role of the parameters in this relationship, we made a simplifying assumption that the synaptic conductance \( g_{\text{syn}}(t) \) rapidly reached its peak (i.e., \( \tau_r \) was small), stayed at this value and started to decay at \( t = t_{\text{peak}} \). In this case \( g(t) = g_{\text{max}} \) on the interval \((0, t_{\text{peak}})\) and the value of \( g_{lo} \) is irrelevant. With this assumption, Equation (5) reduced to

\[ \phi_{LP} = \frac{\tau_s}{P} \ln \frac{g_{\text{max}}}{g^*} + \frac{t_{\text{peak}}}{P}. \]  

(6)

Substituting \( t_{\text{peak}} = \Delta_{\text{peak}} \cdot T_{\text{act}} \) in Equation (6), gave

\[ \phi_{LP} = F(P, g_{\text{max}}, \Delta_{\text{peak}}) = \frac{1}{P} \left( \tau_s \ln \frac{g_{\text{max}}}{g^*} + \Delta_{\text{peak}} T_{\text{act}} \right), \]  

(7)

which we used to describe the LP phase in the C-Dur case. To describe the C-DC case, after substituting \( t_{\text{peak}} = \Delta_{\text{peak}} \cdot DC \cdot P \), we obtained

\[ \phi_{LP} = F(P, g_{\text{max}}, \Delta_{\text{peak}}) = \frac{1}{P} \left( \tau_s \ln \frac{g_{\text{max}}}{g^*} \right) + \Delta_{\text{peak}} \cdot DC. \]  

(8)

Note that these equations also describe the relationship between \( \phi_{LP} \) with \( T_{\text{act}} \) in the C-Dur case, and \( DC \) in the C-DC case).

Equations (7) and (8) can be used to approximate a range of parameters over which \( \phi_{LP} \) is maintained at a constant value \( \phi_c \). To do so, we assumed a specific parameter set, say \( (\hat{P}, \hat{g}_{\text{max}}, \hat{\Delta}_{\text{peak}}) \), satisfies

\[ F(\hat{P}, \hat{g}_{\text{max}}, \hat{\Delta}_{\text{peak}}) = \phi_c, \]

for some fixed phase value, \( \phi_c \). We could now ask whether there are nearby parameters for which phase remains constant, i.e., \( F \) remains equal to \( \phi_c \). The Implicit Function Theorem (Krantz and Parks, 2012) guarantees that this is the case, provided certain derivatives evaluated at \( (\hat{P}, \hat{g}_{\text{max}}, \hat{\Delta}_{\text{peak}}) \) are non-zero, which turns out to be true.
over a large range of parameters. Since the partial derivative with respect to \( \Delta_{\text{peak}} \) of \( F(P, g_{\text{max}}, \Delta_{\text{peak}}) \) at this point is a non-zero constant equal to \( T_{\text{act}}/P \) (or DC) in the C-Dur (or C-DC) case, there is a function \( \Delta_{\text{peak}} = h(P, g_{\text{max}}) \) such that

\[
F(P, g_{\text{max}}, h(P, g_{\text{max}})) = \phi_c
\]

for values of \( P \) and \( g_{\text{max}} \) near \( \hat{P}, \hat{g}_{\text{max}} \). In other words, the Implicit Function Theorem guarantees that small changes in \( P \) and \( g_{\text{max}} \) can be compensated for by an appropriate choice of \( \Delta_{\text{peak}} \) in order to maintain a constant LP phase. A similar analysis can be done by solving for \( g_{\text{max}} \) in terms of \( P \) and \( \Delta_{\text{peak}} \) or by solving for \( P \) in terms of \( g_{\text{max}} \) and \( \Delta_{\text{peak}} \).

Keeping \( g_{\text{max}} \) (respectively, \( \Delta_{\text{peak}} \)) constant in these equations allows us to obtain a relationship between \( P \) and \( \Delta_{\text{peak}} \) (respectively, \( g_{\text{max}} \)), for which \( \phi_{\text{LP}} \) is kept constant at \( \phi_c \). Consider Equations (7) and (8) for fixed values of both \( \phi_{\text{LP}} (= \phi_c) \) and \( g_{\text{max}} \). Then these equations reduce to simple functional relationships where \( \Delta_{\text{peak}} \) can be expressed as a function of \( P \). In the C-DC case, for example, evaluating \( \Delta_{\text{peak}} \) from Equation (8) produces

\[
g_{\text{max}} = g^* \cdot \exp \left( \frac{P}{\tau_s} (\phi_c - \Delta_{\text{peak}}) \right)
\]

(10)

Alternatively, \( g_{\text{max}} \) can be expressed as a function of \( P \). In the C-DC case, evaluating \( \Delta_{\text{peak}} \) from Equation (8) produces

\[
\Delta_{\text{peak}} = \frac{\phi_c}{DC} - \frac{\tau_s}{DC \cdot P} \ln \frac{g_{\text{max}}}{g^*},
\]

(11)

Equation (11) can be used to calculate the range of \( P \) values over which changing \( \Delta_{\text{peak}} \) (from 0 to 1) can maintain a constant phase \( \phi_c \). Solving \( 0 < \Delta_{\text{peak}} < 1 \) using Equation (11) yields

\[
\frac{\tau_s}{\phi_c} \ln \frac{g_{\text{max}}}{g^*} < P_{\text{DC}} < \frac{\tau_s}{\phi_c - DC} \ln \frac{g_{\text{max}}}{g^*}
\]

(12)
Performing the same procedure in the C-Dur case, we find

\[
\frac{\tau_s}{\varphi_c} \ln \frac{g_{\text{max}}}{g^*} < P_{\text{dur}} < \frac{T_{\text{act}}}{\varphi_c} + \frac{\tau_s}{\varphi_c} \ln \frac{g_{\text{max}}}{g^*}. \tag{13}
\]

The lower limits of the two cases \(P_{\text{DC}}\) and \(P_{\text{Dur}}\) are the same. The upper limit for \(P_{\text{DC}}\) is larger than that of \(P_{\text{Dur}}\) if

\[
\varphi_c < DC \left(1 + \frac{\tau_s}{T_{\text{act}}} \ln \frac{g_{\text{max}}}{g^*}\right). \tag{14}
\]

If \(\Delta P\) denotes the range of \(P\) values that respectively satisfy Equation (12) or (13), then \(\Delta P_{\text{DC}} > \Delta P_{\text{Dur}}\) if the inequality given by (14) holds, which it does for \(\tau_s\) and \(g_{\text{max}}\) large enough.

Adding synaptic depression to the model of synaptic dynamics

In a previous modeling study, we explored how the phase of a follower neuron was affected when the inhibitory synapse from an oscillatory neuron to this follower had short-term synaptic depression (Manor et al., 2003). In that study the role of the parameter \(\Delta_{\text{peak}}\) was not considered. It is straightforward to add synaptic depression to Equations (7) and (8) and therefore examine how phase is affected if \(\Delta_{\text{peak}}\) increases with \(P\) and synaptic strength also changes with \(P\) according to the rules of synaptic depression. We will restrict this section to the C-DC case. A similar derivation can be made for the C-Dur case.

An ad hoc model of synaptic depression can be made using a single variable \(s_d\) which will be a periodic function that denotes the extent of depression and takes on values between 0 and 1 (Bose et al., 2004). \(s_d\) decays during the AB/PD burst (from time 0 to \(T_{\text{act}}\), indicating depression) and then recovers during the inter-burst interval (from \(T_{\text{act}}\) to \(P\), indicating recovery). Thus, \(s_d\) can be described by an equation of the form:

\[
\frac{ds_d}{dt} = \begin{cases} -s_d / \tau_\beta & t \mod P \leq T_{\text{act}} \\ (1 - s_d) / \tau_\alpha & T_{\text{act}} < t \mod P < P \end{cases}
\]

Using periodicity, it is straightforward to show that the maximum value of \(s_d\) which occurs at the start of the AB/PD burst, is given by:
Note that $s_{\text{max}}$ is a monotonically increasing function with values between 0 and 1. Its value approaches 1 as $P$ increases, indicating that the synapse becomes stronger. For a complete derivation and description, see (Bose et al., 2004). The effect of synaptic depression on synaptic strength can be obtained by setting

$$g_{\text{max}} = \bar{g}_{\text{max}} \cdot s_{\text{max}}(P)$$

where $s_{\text{max}}$ is given by Equation (15).

**Software, analysis and statistics**

Data were analyzed using MATLAB scripts to calculate the time of burst onset and the phase. Statistical analysis was performed using Sigmaplot 12.0 (Systat). Significance was evaluated with an $\alpha$ value of 0.05, error bars and error values reported denote standard error of the mean (SEM).

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**Competing interests**

The authors declare no competing financial interests.

**References**

Bargmann CI (2012) Beyond the connectome: how neuromodulators shape neural circuits. Bioessays 34:458-465.
Belluscio MA, Mizuseki K, Schmidt R, Kempter R, Buzsaki G (2012) Cross-frequency phase-phase coupling between theta and gamma oscillations in the hippocampus. J Neurosci 32:423-435.
Bidaye SS, Bockemuhl T, Buschges A (2018) Six-legged walking in insects: how CPGs, peripheral feedback, and descending signals generate coordinated and adaptive motor rhythms. J Neurophysiol 119:459-475.
Bose A, Manor Y, Nadim F (2004) The activity phase of postsynaptic neurons in a simplified rhythmic network. J Comput Neurosci 17:245-261.
Brezina V (2010) Beyond the wiring diagram: signalling through complex neuromodulator networks. Philos Trans R Soc Lond B Biol Sci 365:2363-2374.

Bucher D, Marder E (2013) SnapShot: Neuromodulation. Cell 155:482-482 e481.

Bucher D, Prinz AA, Marder E (2005) Animal-to-animal variability in motor pattern production in adults and during growth. J Neurosci 25:1611-1619.

Bucher D, Haspel G, Golowasch J, Nadim F (2015) Central Pattern Generators. eLS.

Buzsaki G, Wang XJ (2012) Mechanisms of gamma oscillations. Annu Rev Neurosci 35:203-225.

Buzsaki G, Tingley D (2018) Space and Time: The Hippocampus as a Sequence Generator. Trends in cognitive sciences 22:853-869.

Cohen AH, Ermentrout GB, Kiemel T, Kopell N, Sigvardt KA, Williams TL (1992) Modeling of Intersegmental Coordination in the Lamprey Central Pattern Generator for Locomotion. Trends in Neurosciences 15:434-438.

Daur N, Nadim F, Bucher D (2016) The complexity of small circuits: the stomatogastric nervous system. Curr Opin Neurobiol 41:1-7.

Dicaprio R, Jordan G, Hampton T (1997) Maintenance of motor pattern phase relationships in the ventilatory system of the crab. J Exp Biol 200:963-974.

Eisen JS, Marder E (1984) A mechanism for production of phase shifts in a pattern generator. J Neurophysiol 51:1375-1393.

Fortune ES, Rose GJ (2001) Short-term synaptic plasticity as a temporal filter. Trends Neurosci 24:381-385.

Goaillard JM, Taylor AL, Schulz DJ, Marder E (2009) Functional consequences of animal-to-animal variation in circuit parameters. Nat Neurosci 12:1424-1430.

Grande LA, Spain WJ (2005) Synaptic depression as a timing device. Physiology (Bethesda, Md) 20:201-210.

Graubard K, Raper JA, Hartline DK (1980) Graded synaptic transmission between spiking neurons. Proc Natl Acad Sci U S A 77:3733-3735.

Greenberg I, Manor Y (2005) Synaptic depression in conjunction with A-current channels promote phase constancy in a rhythmic network. J Neurophysiol 93:656-677.

Grillner S (2006) Biological pattern generation: the cellular and computational logic of networks in motion. Neuron 52:751-766.

Grillner S, El Manira A (2015) The intrinsic operation of the networks that make us locomote. Curr Opin Neurobiol 31:244-249.

Gruhn M, Guckenheimer J, Land B, Harris-Warrick RM (2005) Dopamine modulation of two delayed rectifier potassium currents in a small neural network. J Neurophysiol 94:2888-2900.

Hajos N, Palhalmi J, Mann EO, Nemeth B, Paulsen O, Freund TF (2004) Spike timing of distinct types of GABAergic interneuron during hippocampal gamma oscillations in vitro. J Neurosci 24:9127-9137.

Harris-Warrick RM (2011) Neuromodulation and flexibility in Central Pattern Generator networks. Curr Opin Neurobiol 21:685-692.

Harris-Warrick RM, Johnson BR (2010) Checks and balances in neuromodulation. Front Behav Neurosci 4.
Harris-Warrick RM, Coniglio LM, Levini RM, Gueron S, Guckenheimer J (1995a) Dopamine modulation of two subthreshold currents produces phase shifts in activity of an identified motoneuron. J Neurophysiol 74:1404-1420.

Harris-Warrick RM, Coniglio LM, Barazangi N, Guckenheimer J, Gueron S (1995b) Dopamine modulation of transient potassium current evokes phase shifts in a central pattern generator network. J Neurosci 15:342-358.

Hasenstaub A, Shu Y, Haider B, Kraushaar U, Duque A, McCormick DA (2005) Inhibitory postsynaptic potentials carry synchronized frequency information in active cortical networks. Neuron 47:423-435.

Hasselmo ME, Bodelon C, Wyble BP (2002) A proposed function for hippocampal theta rhythm: separate phases of encoding and retrieval enhance reversal of prior learning. Neural Comput 14:793-817.

Hooper SL (1997a) Phase maintenance in the pyloric pattern of the lobster (Panulirus interruptus) stomatogastric ganglion. J Comput Neurosci 4:191-205.

Hooper SL (1997b) The pyloric pattern of the lobster (Panulirus interruptus) stomatogastric ganglion comprises two phase-maintaining subsets. J Comput Neurosci 4:207-219.

Hooper SL (1998) Transduction of temporal patterns by single neurons. Nat Neurosci 1:720-726.

Hooper SL, Buchman E, Weaver AL, Thuma JB, Hobbs KH (2009) Slow conductances could underlie intrinsic phase-maintaining properties of isolated lobster (Panulirus interruptus) pyloric neurons. J Neurosci 29:1834-1845.

Johnson BR, Kloppenburg P, Harris-Warrick RM (2003) Dopamine modulation of calcium currents in pyloric neurons of the lobster stomatogastric ganglion. J Neurophysiol 90:631-643.

Johnson BR, Schneider LR, Nadim F, Harris-Warrick RM (2005) Dopamine modulation of phasing of activity in a rhythmic motor network: contribution of synaptic and intrinsic modulatory actions. J Neurophysiol 94:3101-3111.

Jordan LM, Slawinska U (2011) Chapter 12--modulation of rhythmic movement: control of coordination. Prog Brain Res 188:181-195.

Katz PS (2016) Evolution of central pattern generators and rhythmic behaviours. Philos Trans R Soc Lond B Biol Sci 371:20150057.

Kayser C, Montemurro MA, Logothetis NK, Panzeri S (2009) Spike-phase coding boosts and stabilizes information carried by spatial and temporal spike patterns. Neuron 61:597-608.

Kiehn O (2016) Decoding the organization of spinal circuits that control locomotion. Nat Rev Neurosci 17:224-238.

Kloppenburg P, Levini RM, Harris-Warrick RM (1999) Dopamine modulates two potassium currents and inhibits the intrinsic firing properties of an identified motor neuron in a central pattern generator network. J Neurophysiol 81:29-38.

Kopell N, Whittington MA, Kramer MA (2011) Neuronal assembly dynamics in the beta1 frequency range permits short-term memory. Proc Natl Acad Sci U S A 108:3779-3784.
Krantz SG, Parks HR (2012) The implicit function theorem: history, theory, and applications: Springer Science & Business Media.

Kvarta MD, Harris-Warrick RM, Johnson BR (2012) Neuromodulator-evoked synaptic metaplasticity within a central pattern generator network. J Neurophysiol 108:2846-2856.

Le Gal JP, Dubuc R, Smarandache-Wellmann C (2017) Coordination of Rhythmic Movements. In: Neurobiology of Motor Control: Fundamental Concepts and New Directions (Hooper SL, Buschges A, eds). Hoboken, New Jersey: Wiley-Blackwell.

Lindner B A brief introduction to some simple stochastic processes. Stochastic Methods in Neuroscience 1.

Luther JA, Robie AA, Yarotsky J, Reina C, Marder E, Golowasch J (2003) Episodic bouts of activity accompany recovery of rhythmic output by a neuromodulator- and activity-deprived adult neural network. J Neurophysiol 90:2720-2730.

Ma C, LaMotte RH (2007) Multiple sites for generation of ectopic spontaneous activity in neurons of the chronically compressed dorsal root ganglion. J Neurosci 27:14059-14068.

Mamiya A, Nadim F (2004) Dynamic interaction of oscillatory neurons coupled with reciprocally inhibitory synapses acts to stabilize the rhythm period. J Neurosci 24:5140-5150.

Mamiya A, Manor Y, Nadim F (2003) Short-term dynamics of a mixed chemical and electrical synapse in a rhythmic network. J Neurosci 23:9557-9564.

Manor Y, Nadim F, Abbott LF, Marder E (1997) Temporal dynamics of graded synaptic transmission in the lobster stomatogastric ganglion. J Neurosci 17:5610-5621.

Manor Y, Bose A, Booth V, Nadim F (2003) Contribution of synaptic depression to phase maintenance in a model rhythmic network. J Neurophysiol 90:3513-3528.

Marder E (2012) Neuromodulation of neuronal circuits: back to the future. Neuron 76:1-11.

Marder E, Bucher D (2001) Central pattern generators and the control of rhythmic movements. Curr Biol 11:R986-996.

Marder E, Goaillard JM (2006) Variability, compensation and homeostasis in neuron and network function. Nat Rev Neurosci 7:563-574.

Marder E, Bucher D (2007) Understanding circuit dynamics using the stomatogastric nervous system of lobsters and crabs. Annu Rev Physiol 69:291-316.

Marder E, O’Leary T, Shruti S (2014) Neuromodulation of circuits with variable parameters: single neurons and small circuits reveal principles of state-dependent and robust neuromodulation. Annu Rev Neurosci 37:329-346.

Marder E, Bucher D, Schulz DJ, Taylor AL (2005) Invertebrate central pattern generation moves along. Curr Biol 15:R685-699.

Markram H, Wang Y, Tsodyks M (1998) Differential signaling via the same axon of neocortical pyramidal neurons. Proc Natl Acad Sci U S A 95:5323-5328.
Masino MA, Fetcho JR (2005) Fictive swimming motor patterns in wild type and mutant larval zebrafish. J Neurophysiol 93:3177-3188.

McLelland D, Paulsen O (2009) Neuronal oscillations and the rate-to-phase transform: mechanism, model and mutual information. J Physiol 587:769-785.

Mouser C, Nadim F, Bose A (2008) Maintaining phase of the crustacean tri-phasic pyloric rhythm. J Math Biol 57:161-181.

Mullins OJ, Hackett JT, Buchanan JT, Friesen WO (2011) Neuronal control of swimming behavior: comparison of vertebrate and invertebrate model systems. Prog Neurobiol 93:244-269.

Nadim F, Manor Y (2000) The role of short-term synaptic dynamics in motor control. Curr Opin Neurobiol 10:683-690.

Nadim F, Bucher D (2014) Neuromodulation of neurons and synapses. Curr Opin Neurobiol 29:48-56.

Nadim F, Booth V, Bose A, Manor Y (2003) Short-term synaptic dynamics promote phase maintenance in multi-phasic rhythms. Neurocomputing 52-4:79-87.

Oren I, Mann EO, Paulsen O, Hajos N (2006) Synaptic currents in anatomically identified CA3 neurons during hippocampal gamma oscillations in vitro. J Neurosci 26:9923-9934.

Panzeri S, Brunel N, Logothetis NK, Kayser C (2010) Sensory neural codes using multiplexed temporal scales. Trends Neurosci 33:111-120.

Peck JH, Gaier E, Stevens E, Repicky S, Harris-Warrick RM (2006) Amine modulation of Ih in a small neural network. J Neurophysiol 96:2931-2940.

Pulver SR, Bayley TG, Taylor AL, Berni J, Bate M, Hedwig B (2015) Imaging fictive locomotor patterns in larval Drosophila. J Neurophysiol 114:2564-2577.

Rabbah P, Nadim F (2007) Distinct synaptic dynamics of heterogeneous pacemaker neurons in an oscillatory network. J Neurophysiol 97:2239-2253.

Rosenbaum P, Marder E (2018) Graded Transmission without Action Potentials Sustains Rhythmic Activity in Some But Not All Modulators That Activate the Same Current. J Neurosci 38:8976-8988.

Schulz DJ, Goaillard JM, Marder E (2006) Variable channel expression in identified single and electrically coupled neurons in different animals. Nat Neurosci 9:356-362.

Skinner FK, Mulloney B (1998) Intersegmental coordination in invertebrates and vertebrates. Curr Neurobiol 8:725-732.

Somogyi P, Klausberger T (2005) Defined types of cortical interneurone structure space and spike timing in the hippocampus. J Physiol 562:9-26.

Soofi W, Goeritz ML, Kispersky TJ, Prinz AA, Marder E, Stein W (2014) Phase maintenance in a rhythmic motor pattern during temperature changes in vivo. J Neurophysiol 111:2603-2613.

Stein PSG (2018) Central pattern generators in the turtle spinal cord: selection among the forms of motor behaviors. J Neurophysiol 119:422-440.

Stein W (2009) Modulation of stomatogastric rhythms. J Comp Physiol A Neuroethol Sens Neural Behav Physiol 195:989-1009.

Tang LS, Taylor AL, Rinberg A, Marder E (2012) Robustness of a rhythmic circuit to short- and long-term temperature changes. J Neurosci 32:10075-10085.
Tohidi V, Nadim F (2009) Membrane resonance in bursting pacemaker neurons of an oscillatory network is correlated with network frequency. J Neurosci 29:6427-6435.

Tseng HA, Nadim F (2010) The membrane potential waveform of bursting pacemaker neurons is a predictor of their preferred frequency and the network cycle frequency. J Neurosci 30:10809-10819.

Tseng HA, Martinez D, Nadim F (2014) The frequency preference of neurons and synapses in a recurrent oscillatory network. J Neurosci 34:12933-12945.

Vidal-Gadea A, Topper S, Young L, Crisp A, Kressin L, Elbel E, Maples T, Brauner M, Erbguth K, Axelrod A, Gottschalk A, Siegel D, Pierce-Shimomura JT (2011) Caenorhabditis elegans selects distinct crawling and swimming gaits via dopamine and serotonin. Proc Natl Acad Sci U S A 108:17504-17509.

Wang XJ (2010) Neurophysiological and computational principles of cortical rhythms in cognition. Physiol Rev 90:1195-1268.

Weimann JM, Meyrand P, Marder E (1991) Neurons that form multiple pattern generators: identification and multiple activity patterns of gastric/pyloric neurons in the crab stomatogastric system. J Neurophysiol 65:111-122.

Wenning A, Hill AA, Calabrese RL (2004) Heartbeat control in leeches. II. Fictive motor pattern. J Neurophysiol 91:397-409.

Wenning A, Norris BJ, Gunay C, Kueh D, Calabrese RL (2018) Output variability across animals and levels in a motor system. eLife 7.

Zhang C, Guy RD, Mulloney B, Zhang Q, Lewis TJ (2014) Neural mechanism of optimal limb coordination in crustacean swimming. Proc Natl Acad Sci U S A 111:13840-13845.

Zhang Y, Bose A, Nadim F (2008) Predicting the activity phase of a follower neuron with A-current in an inhibitory network. Biol Cybern 99:171-184.

Zhang Y, Bose A, Nadim F (2009) The influence of the A-current on the dynamics of an oscillator-follower inhibitory network. SIAM J Appl Dyn Syst 8:1564-1590.

Zhao S, Sheibanie AF, Oh M, Rabbah P, Nadim F (2011) Peptide neuromodulation of synaptic dynamics in an oscillatory network. J Neurosci 31:13991-14004.
A1. Schematic diagram showing that a follower neuron (F) strongly inhibited by a bursting oscillatory neuron (O) with period $P$ can produce rebound bursts with the same period at a latency $\Delta t$. A2. If the period of O changes to a new value ($P_2$), the new F burst latency ($\Delta t_2$) typically falls between two extremes: it could stay constant (top trace) or change proportionally to $P_2$, so that the burst phase ($\Delta t / P$) remains constant (middle trace). B. Example traces of the pyloric pacemaker PD neuron and the follower LP neuron represent the O and F relationship in panel A. Here, the PD neuron is voltage clamped and a pre-recorded waveform of the same neuron is used to drive this neuron to follow different cycle periods. The LP neuron follows the same period because of the synaptic input it receives. C. A measurement of the LP neuron burst onset time ($\Delta t$) with respect to the onset of the PD neuron burst shows that $\Delta t$ falls between the two limits of constant latency and constant phase. Dotted curves represent constant latency matched to the latencies at the two extreme $P$ values.
Figure 2: Inputs to the LP neuron influence burst time, spike number and interval

A. Simultaneous intracellular recording of the LP neuron and extracellular recording of the lateral ventricular nerve (lvn), containing the axons of the LP, PD and PY neurons (arrows). Period (P) and the burst onset time (Δt) of the LP neuron are defined in reference to the pacemaker group (PD) burst. B. Blocking the AB and PY synaptic inputs (10 µM picrotoxin) to the LP neuron disrupts its bursting oscillations. C. The LP neuron, in picrotoxin, was driven with a noise current input (INoise) for 60 minutes. In response, the LP neuron produced an irregular pattern of bursting. Specific inter-burst intervals (IBIs) were tagged and used for burst-triggered averaging. D. Example of burst-trigger-averaged input current (IBTA, green). Individual traces are shown in grey. E. For each IBI (300, 500, 700, 900 ms), IBTA was calculated and normalized to the (negative) peak value of IBTA for IBI=300 ms. Different traces in each panel show the IBTA of different preparations. F. The mean (across preparations) of the normalized IBTAS shown in panel E. G. Traces in panel F normalized by IBI. H-K. Four parameters define the shape of the IBTA: peak amplitude Iamp (H), peak phase Δpeak (I), slopeup (J) and slopedown (K) across preparations. IBI had a significant effect on amplitude Iamp (p<0.001), peak phase Δpeak (p<0.001), slopeup (p<0.001) and slopedown (p=0.002).
Figure 3: Cycle period and synaptic strength affect the phase of LP burst onset in opposite directions.

A. The synaptic input to the LP neuron was measured by voltage clamping it at a holding potential of -50mV during ongoing oscillations. The onset of the pacemaker (AB/PD) activity is seen as a kink in the synaptic current (ILP, blue). Dashed line: 0 nA.

B. Synaptic input averaged across (last 5 of 30) cycles from 9 different LP neurons. Traces are aligned to the onset of the PD neuron burst (dotted vertical red line; see panel A), normalized by the cycle period and terminated at the end of the downslope (coincident with the first LP action potential when present). The blue trace shows the average.

C. An example of the LP neuron driven by the realistic synaptic waveform in dynamic clamp. The burst onset time (Δt) was measured relative to the AB/PD onset and used to measure the LP phase (φLP). gmax denotes the conductance amplitude.

D. Mean φLP (N=9 preparations) shown as a function of P and fit with the function given by Equation (8) (fit values τs = 26.0 ms, g* = 0.021 µS and Δpeak·DC = 0.43).

E. Mean φLP plotted against gmax also shown with the fit to Equation (8).

F. Heat map, obtained from fitting Equation (8) to the data in panels D and E, shows φLP as a function of both gmax and P. Black curves show the level sets of phase constancy for three values of φLP (0.47, 0.49, and 0.52).
Figure 4: The constant duty cycle of synaptic conductance is a major factor in phase maintenance.

A. The change in $\phi_{LP}$ values with $P$ are compared with the constant phase (solid curve) and constant latency (dashed curve) extremes. Lime traces show the usual values of $\phi_{LP}$, calculated from the LP burst onset latency with respect to the onset of the PD burst. Lavender traces show $\phi_{LP}$ calculated from the LP burst onset latency with respect to the end of the PD burst. Data shown are the same as in Fig. 3D for $g_{max}=0.4 \mu S$.

B. Schematic diagram shows the latency of LP burst onset measured with respect to the (estimated) onset and end of the PD burst in the dynamic clamp experiments (see Methods). Bottom panel shows the synaptic current waveform measured in the voltage-clamped LP neuron during ongoing pyloric activity. Top panel shows the dynamic clamp injection of the synaptic conductance waveform into the LP neuron. The current waveform of the bottom panel is aligned to the conductance waveform of the top panel for the comparison used in determining the PD burst onset and end in the top panel.
Figure 5: Four parameters describing synaptic shape were varied in the experimental paradigm.

A. A triangle shaped conductance was used to mimic the synaptic input to the LP neuron. B. The triangular waveform can be described by period ($P$), duration ($T_{act}$), peak time ($t_{peak}$) and amplitude ($g_{max}$). C. In dynamic clamp runs, the synapse duration $T_{act}$ was kept constant at 300 ms (C-Dur) or maintained at a constant duty cycle ($T_{act} / P$) of 0.3 (C-DC) across all values of $P$. D. Intracellular voltage recording of the LP neuron during a dynamic clamp stimulation run using the triangle conductance (in picrotoxin). The burst onset time ($\Delta t$, calculated in reference to the synaptic conductance onset) was used to calculate the activity phase ($\phi_{LP} = \Delta t / P$).
Figure 6: The LP burst onset phase decreases as a function of P, but increases as a function of $g_{\text{max}}$ and $\Delta_{\text{peak}}$.

Periodic injection of an inhibitory triangular waveform conductance into the LP neuron (in picrotoxin) produced bursting activity from which $\phi_{\text{LP}}$ was calculated. The parameters $g_{\text{max}}, \Delta_{\text{peak}}$ and $P$ were varied across runs for both C-Dur and C-DC cases.

A. $\phi_{\text{LP}}$ decreases as a function of $P$. A1. Intracellular recording of an LP neuron showing a C-DC conductance input across five periods. A2. $\phi_{\text{LP}}$ for the example shown in A1 plotted as a function of $P$ (for $g_{\text{max}} = 0.4 \, \mu\text{S}, \Delta_{\text{peak}} = 0.5$) for both C-Dur and C-DC cases. $\phi_{\text{LP}}$ decreases rapidly with $P$ and the drop is larger for the C-Dur case. A3. $\phi_{\text{LP}}$ decreased with $P$ in both the C-DC case (Three-Way RM ANOVA, $p<0.001, F=100.7$) and the C-Dur case (Three-Way RM ANOVA, $p<0.001, F=466.4$) for all values of $\Delta_{\text{peak}}$. The range of $\phi_{\text{LP}}$ drop was greater for the C-Dur case compared to the C-DC case.

B. $\phi_{\text{LP}}$ increases as a function of $g_{\text{max}}$. B1. Intracellular recording of an LP neuron showing the conductance input across three values of $g_{\text{max}}$. B2. $\phi_{\text{LP}}$ for the example shown in B1 plotted as a function of $P$ (for $P = 500 \, \text{ms}, \Delta_{\text{peak}} = 0.25$) shows a small increase for both C-Dur and C-DC cases. B3. $\phi_{\text{LP}}$ increased with $g_{\text{max}}$ in almost all trials for both C-DC and C-Dur cases and all values of $\Delta_{\text{peak}}$. C. $\phi_{\text{LP}}$ increases as a function of $\Delta_{\text{peak}}$. C1. Intracellular recording of the LP neuron showing the conductance input for five values of $\Delta_{\text{peak}}$ C2. $\phi_{\text{LP}}$ for the example neuron in C1 plotted as a function of $\Delta_{\text{peak}}$ (for $P = 500 \, \text{ms}, g_{\text{max}} = 0.4 \, \mu\text{S}$) for both C-DC and C-Dur cases. C3. $\phi_{\text{LP}}$ increased with $\Delta_{\text{peak}}$ for both C-DC and C-Dur cases and for all values of $g_{\text{max}}$. In all panels, error bars show standard deviation.
Figure 7: Sensitivity analysis shows that $\phi_{LP}$ increases more effectively if $g_{\text{max}}$ and $\Delta_{\text{peak}}$ increase together.

A. The sensitivity of $\phi_{LP}$ to local changes in $g_{\text{max}}$ and $\Delta_{\text{peak}}$ was averaged across all values of $P$ for the C-DC case. Sensitivity was largest if both parameters were increased together ($g_{\text{max}} + \Delta_{\text{peak}}$) and smallest if they were varied in opposite directions ($g_{\text{max}} - \Delta_{\text{peak}}$; One-Way RM-ANOVA, p<0.001, F=3.330).

B. The same sensitivity analysis in the C-Dur case shows similar results (One-Way RM-ANOVA, p<0.001, F=2.892). In both panels, error bars show standard deviation.
Figure 8: Simultaneous increase of both Δ\text{peak} and g_{max} across their range of values can produce phase maintenance across a large P range in the C-DC case and a much smaller P range in the C-Dur case.

A. Heat map plots of the function Φ (see Methods), plotted for the range of values of P and Δ\text{peak} and 4 values of g_{max} for the C-DC (A1) and C-Dur (A2) cases. The white curves show the level set of φ_{LP} = 0.34, shown as an example of phase constancy. The color maps are interpolated from sampled data (see Methods; N=9 experiments). The locations of the sampled data are marked by black dots. B. Heat map for the level sets φ_{LP} = 0.34 for the C-DC (B1) and C-Dur (B2) cases. Range of colors in each panel indicate the range of P values for which φ_{LP} could remain constant at 0.34 for each case, as indicated by the grey arrows on the side of the heatmap color legend. C. The range (ΔP) of P values for which φ_{LP} could remain constant at any value between 0.2 and 0.8 for the C-DC (C1) and C-Dur cases (C2). Filled circles show the values shown in panel B. The LP neuron cannot achieve φ_{LP} values below 0.3 in the C-DC case. For φ_{LP} values between 0.3 and ~0.65, the range was larger in C-DC case.
Figure 9: Model prediction of the range of phase constancy.

A. For the C-DC case, a constant phase of $\phi_{LP} = 0.34$ can be maintained across a range of cycle periods $P$ when $g_{max}$ is constant (at 335 nS; blue plane) and $\Delta_{peak}$ varies from 0 to 1 according to Equation (11) (blue), or when $\Delta_{peak}$ is fixed (at 0.5; green plane) and $g_{max}$ varies from 200 to 800 nS according to Equation (10). Alternatively, $g_{max}$ and $\Delta_{peak}$ can covary to maintain phase, as in a depressing synapse, where $g_{max}$ varies with $P$ according to Equation (16), and $\Delta_{peak}$ is calculated for each $P$ and $g_{max}$ value according to Equation (11). As seen in the 2D coordinate-plane projections of the 3D graph (right three graphs), the range of $P$ values for which phase constancy is achieved is largest when $g_{max}$ and $\Delta_{peak}$ covary (dotted lines show limits of $P$ for phase constancy). The depressing synapse conductance is chosen to be 335 nS at $P = 1$ s.

B, C. A comparison between the C-DC and C-Dur cases shows that in the latter case a constant phase of $\phi_{LP} = 0.34$ can be maintained across a larger range of $P$ values when $\Delta_{peak}$ increases with $P$ (and $g_{max}$ is fixed at 400 nS) according to Equation (11). The relationship of $\Delta_{peak}$ and $P$ is shown in B for $\phi_{LP} = 0.34$. C shows the range of $P$ values ($\Delta P$) of cycle periods for which phase remains constant at $\phi_{LP}$. If $g_{max}$ also varies with $P$, as in a depressing synapse (red; Equation (16)), the range of $P$ values for which phase is constant is further increased. (Dotted line: $\phi_{LP} = 0.34$.)
Source Data Files

**Figure 2-1. File: Figure2_sourcedata.xlsx**

This Excel file contains 4 sheets, including all measured attributes of the burst-triggered average current (IBTA) for different IBIs (N=23) as shown in Fig. 2H-2K.

**Figure 7-1. File: Figure7_sourcedata.xlsx**

This Excel file contains 2 sheets for the C-DC and C-Dur cases. These sheets include all sensitivity values for each value of $P$, at each $g_{\text{max}}$ and each $\Delta_{\text{peak}}$ in all 8 directions: $(+g_{\text{max}}, +\Delta_{\text{peak}}, -g_{\text{max}}, -\Delta_{\text{peak}}, +g_{\text{max}} \& +\Delta_{\text{peak}}, -g_{\text{max}} \& -\Delta_{\text{peak}}, +g_{\text{max}} \& -\Delta_{\text{peak}}, +g_{\text{max}} \& +\Delta_{\text{peak}})$. Fig. 7 shows the sensitivities, averaged across all $P$ values, and averaged across aligned directions: $[+g_{\text{max}} \& -g_{\text{max}}]; [+\Delta_{\text{peak}} \& -\Delta_{\text{peak}}]; [+g_{\text{max}} \& +\Delta_{\text{peak}} \& -g_{\text{max}} \& -\Delta_{\text{peak}}]; [+g_{\text{max}} \& -\Delta_{\text{peak}} \& +g_{\text{max}} \& -\Delta_{\text{peak}}].$