X-ray colour-colour selection for heavily absorbed AGN

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ABSTRACT

We present a method for the identification of heavily absorbed AGN ($N_{\text{H}} > 10^{23}$ cm$^{-2}$) from X-ray photometric data. We do this using a set of XMM-Newton reference spectra of local galaxies for which we have accurate $N_{\text{H}}$ information, as described in Brightman & Nandra. The technique uses two rest-frame hardness ratios which are optimised for this task, which we designate HR1 (2-4/1-2 keV) and HR2 (4-16/2-4 keV). The selection method exploits the fact that while obscured AGN appear hard in HR2 due to absorption of the intrinsic source flux below $\sim 4$ keV, they appear soft in HR1 due to excess emission originating from scattered source light, thermal emission, or host galaxy emission. Such emission is ubiquitous in low redshift samples. The technique offers a very simple and straightforward way of estimating the fraction of obscured AGN in samples with relatively low signal-to-noise ratio in the X-ray band. We apply this technique to a moderate redshift ($z \sim 1$) sample of AGN from the Chandra Deep Field North, finding that 61% of this sample has $N_{\text{H}} > 10^{23}$ cm$^{-2}$. A clear and robust conclusion from our analysis, is that in deep surveys the vast majority of sources do not show hardness ratios consistent with a simple absorbed power law. The ubiquity of complex spectra in turn shows that simple hardness ratio analysis will not yield reliable obscuration estimates, justifying the more complex colour-colour analysis described in this paper. While this method does very well at separating sources with $N_{\text{H}} > 10^{23}$ cm$^{-2}$ from sources with lower $N_{\text{H}}$, only X-ray spectroscopy can identify Compton thick sources, through the detection of the Fe Kα line. This will be made possible with the high throughput X-ray spectral capabilities of ATHENA.

Key words: galaxies: active - galaxies: Seyfert - X-rays: galaxies

1 INTRODUCTION

The nature and amount of obscuration in active galactic nuclei (AGN) is an important issue for a range of astrophysical questions, such as unification of AGN (eg. Antonucci 1993) and the role of AGN in galaxy evolution (Hopkins et al. 2006). Heavily absorbed AGN can be identified using their X-ray spectra, via direct measurement of the line of sight column density. At the highest column densities ($N_{\text{H}} > 10^{24}$ cm$^{-2}$, where even the direct X-rays are suppressed, obscured AGN can still be identified via the presence of X-ray reflection, with the spectrum consisting of a flat X-ray continuum accompanied by a high equivalent width Fe Kα line (Matt, Brandt & Fabian 1996). The measurement of column densities or identification of reflection features requires good quality X-ray spectra with enough counts to carry out spectral fitting and to reject alternative hypotheses. This is challenging, even in the local universe, as the flux suppression in the 2-10 keV band is significant once the $N_{\text{H}}$ reaches a few $10^{23}$ cm$^{-2}$ (Brightman & Nandra 2011a). Many alternative methods of identifying heavily absorbed sources, where the X-ray spectrum itself is not good enough to do so, have been presented in the literature. These often use the ratio of the observed X-ray flux to the flux at some wavelength at which the effect of the obscuration is assumed to be small, and which should therefore be representative of the AGN bolometric flux. Examples of these assumed “isotropic indicators” are the MIR, [O III] line or [Ne V] line (Diamond-Stanic et al. 2009; LaMassa et al. 2010; Gilli et al. 2010). Obscured AGN may be found via weak observed X-ray emission with respect to the isotropic indicator, with the X-ray flux suppression attributed to absorption. A popular technique which uses this method was presented by Bassani et al. (1999) and uses $F_X/F_{[\text{OIII}]}$ in conjunction with the Fe Kα line equivalent width as proxies for $N_{\text{H}}$, though $F_X/F_{[\text{OIII}]}$ has henceforth often been used alone. The use of isotropic indicators are complicated when the isotropic indicator has contributions from non nuclear sources (eg. star-formation processes con-
tribute to both MIR and [O III] fluxes), or where the source is intrinsically X-ray weak. Furthermore, techniques have even been developed that use data entirely outside of the X-ray band to detect X-ray obscuration (e.g. Daddi et al. 2007; Fiore et al. 2008). Arguably though, the best place to judge X-ray absorption, is in the X-ray band alone. This is sometimes done using ‘hardness ratios’, which is a model independent technique, and utilises X-ray photometry. Traditionally, hardness ratios are calculated using observed counts, often in the standard soft (0.5-2.0 keV) and hard (2.0-10.0 keV) bands. Sources with large hardness ratio are assumed absorbed, while sources with low ratios assumed unabsorbed. However this method is very crude and may not exclusively separate the obscured sources from the unobscured sources if spectral complexity is present. At other wavelengths, colour-colour selection is an effective method of identifying sources populations based on photometry. Here, we seek a novel approach to the identification of X-ray absorption from X-ray photometric data, using colour-colour selection, formed from X-ray hardness ratios. While work using a combination of hardness ratios to infer absorption has been presented previously, for example results by Della Ceca et al. (1999) using ASCA data, we seek to improve on these using knowledge from local samples with data of high spectral quality. This better understanding of hardness ratios can then be applied in cases with much poorer data quality, such as in deep surveys of the distant Universe, and help elucidate the observed relationships seen between obscuration, luminosity and redshift (e.g. Ueda et al. 2003; Haslinger 2008; Brightman & Ueda, in preparation).

2 X-RAY DATA

For the task of finding an effective colour-colour selection technique for obscured AGN, we utilise a set of 126 XMM-Newton spectra of local galaxies, for which we have carried out a detailed spectral analysis of in Brightman & Nandra (2011a, BN11). The parent sample of this set was the extended IRAS 12 micron galaxy sample of Rush et al. (1993). The mid-IR selection of this sample should ensure relatively little bias against obscured objects (Spinoglio & Malkan 1989). The X-ray spectra of these galaxies were fitted in the standard soft (0.5-2.0 keV) and hard (2.0-10.0 keV) bands. Sources with large hardness ratio are assumed absorbed, while sources with low ratios assumed unabsorbed.

We therefore plot two proposed selection criteria. First (1), the solid wedge selects all heavily absorbed sources (N_H > 10^{24} cm^{-2}). As the sample is local (z < 0.1), we treat the X-ray spectral data as rest-frame. The sample consists of all galaxy types, from Seyfert galaxies to H II and LINER galaxies, with a wide range of L_x, N_H and a high fraction of absorbed and Compton thick sources. For a more detailed description of the X-ray spectral analysis, the X-ray spectral properties of the sample, and the new Monte-Carlo models, see BN11. For the optical activity classification and 12 µm data, see Brightman & Nandra (2011b).

3 COLOUR-COLOUR SELECTION FOR HEAVILY OBSCURED AGN

In Fig 2, we present a colour-colour plot which uses our hardness ratios using our reference spectra data. On it, we also show the N_H of each source, represented by colour. There is a cluster of sources at HR1 ≃ HR2 = 0.0. These are objects which show little absorption above 1 keV. Both ratios harden as the N_H increases until ~10^{23} cm^{-2}. At this point, the soft-excess emission begins to dominate below 4 keV and hence HR1 softens, while HR2 remains hard. It can be seen that the heavily absorbed sources (N_H > 10^{23} cm^{-2}) are separated well from the others using these hardness ratios in combination, being hard in HR2, while soft in HR1. We therefore plot two proposed selection criteria. First (1), the dotted line selects almost all (38/41=93%) heavily absorbed sources, with only a few (3/41=7%) contaminants with a lower N_H. Secoundly (2), the solid wedge selects all (16/16=100%) Compton thick sources. This region is contaminated by 20/36 (56%) sources with N_H < 10^{23} cm^{-2}, though only by 3/36 (8%) sources with N_H < 10^{22} cm^{-2}.

The selection criteria are as follows. For N_H > 10^{23} cm^{-2},

(1) HR2 > 0.62 × HR1 + 0.38 and for Compton thick AGN

(2) HR2 > 0.25 AND HR2 < 2.28 × HR1 + 0.36

Calculation of the rest-frame hardness ratios is in practice very straightforward. For example, counts or count rates can be converted to flux to be used in this selection.
Figure 2. Our colour-colour plot using the hardness ratios HR1 (2-4/1-2 keV) and HR2 (4-16/2-4 keV), calculated from model photon fluxes extracted from our set of 126 XMM-Newton reference spectra. The colour of each data point corresponds to the value of the measured $N_{\text{H}}$ in that spectrum. The dotted line and the solid wedge are proposed selection regions for heavily obscured AGN, as described in the text.

4 DISCUSSION AND SUMMARY

In this paper we have aimed to identify an effective colour-colour selection for obscured AGN from X-ray data and improve upon the use of the classical hardness ratio. In Fig. 4 we plot histograms of the classical HR (with photon flux, rather than counts) for unabsorbed sources, absorbed Compton thin sources and Compton thick sources from Brightman & Nandra (2011a). This confirms our view that the hardness ratio is not an effective measure of obscuration, as all Compton thick sources appear in fact soft, while both Compton thin and unabsorbed sources show the full range of hardness ratio.

We have shown that colour-colour selection is able to break these degeneracies. The choice of bands for our colour-colour selection stems from the observation that while obscured AGN do appear hard above $\sim 4$ keV, they are usually dominated by a soft excess component below 4 keV. This component is often scattered nuclear light, which is believed to have undergone Thompson scattering off hot electrons far from the nucleus, and hence out-with the obscuring material. It typically represents only a few percent of the intrinsic emission (Turner et al. 1997). Soft X-ray emission in obscured AGN can also be attributed to circumnuclear gas photoionised by the AGN (Guainazzi & Bianchi 2007), or to star-formation related processes in the host galaxy, such as X-ray binaries. In all cases, this emission is assumed to be not subjected to the obscuration of AGN, and thus appears soft in HR1.

Della Ceca et al. (1999) presented an analysis of hard X-ray selected sources from ASCA utilising very technique using the HEASARC Portable Interactive Multimission Software (pimms). This requires the assumption of a model. For deep X-ray fields, the use of a $\Gamma = 1.4$ power-law with no absorption, roughly consistent with the cosmic X-ray background, is appropriate.

We can then use this new identification technique to make inferences about the heavily obscured AGN population from currently available X-ray survey data. Here we apply the scheme to the Chandra Deep Field North survey (CDFN). Alexander et al. (2003) present photometric data from this survey in the 0.5-1 keV, 1-2 keV and 2-8 keV bands. These correspond exactly to our band selection for $z = 1$ sources. Redshift data for this catalogue have been provided by Trouille et al. (2008), including photometric redshifts. Using these, we select sources with $0.8 < z < 1.2$. The hardness ratio plot for this survey is presented in Fig. 3. Here we also show tracks of a simple absorbed power-law, calculated using the $\text{zwabs}*\text{power-law}$ models in XSPEC with $\Gamma = 1.7, 1.9$ and 2.1, from $N_{\text{H}} = 10^{20}-10^{24}$ cm$^{-2}$.

We find that for the 107 $0.8 < z < 1.2$ CDFN sources we have selected, 65 (61%) lie above the selection line. We can also use these hardness ratios to infer information about the spectral complexity present in these sources. Only 39 (36%) of the CDFN sources we study here show spectra consistent with a simple absorbed power-law, where the majority of sources present more complex spectra, due to soft excess emission.

1 an online calculator can be found here: https://heasarc.gsfc.nasa.gov/Tools/w3pimms.html
similar hardness ratios to ours. They also use these to infer absorption and spectral complexity in their sample. We have improved upon this technique to provide specific selection criteria for AGN with \( N_H > 10^{23} \text{ cm}^{-2} \) as well as Compton thick AGN, and to extend the analysis to higher redshifts. Furthermore, Noguchi et al. (2009) use similar hardness ratios in their study of the scattered fraction in AGN, also using data from 2XMM. They use the hardness ratios to select AGN with very small scattered fractions, in the search for buried AGN first described in Ueda et al. (2007). They also note, however, that heavily obscured AGN can be selected well using these hardness ratios, though this is not the focus of their work.

Our results here provide a valuable technique for the identification of obscured AGN from photometry alone, when a spectral fit is not possible or badly constrained, and provides a direct method for investigating the nature of obscuration in AGN. As we have shown, it has applications for deep field, high redshift studies, where there are few good quality spectra, though above \( z \sim 1 \), the 1-2 keV restframe band begins to be redshifted out of most X-ray telescope bandpasses. The redshift of the source should be known also, for correct application of the narrow bands. As Noguchi et al. (2009) suggested, it is well suited for application to wider and shallower surveys, such as the \( \text{XMM-Newton} \) serendipitous source catalogue, especially as the photometry provided in that catalogue matches well to our choice of bands. While this method does very well at separating sources with \( N_H > 10^{23} \text{ cm}^{-2} \) from sources with lower \( N_H \), only X-ray spectroscopy can identify Compton thick sources, through the detection of the Fe K\( \alpha \) line. This will be made possible for the low fluxes considered here with the high throughput X-ray spectral capabilities of \( \text{ATHENA} \).
We have also used our hardness ratio analysis to make inferences about the spectral complexity CDFN sources as Della Ceca et al. (1999) did for ASCA sources. We find that a significant number of CDFN sources at $z \sim 1$ show spectral shapes more complicated than a simple absorbed power-law due to soft-excess emission. This has implications for the X-ray selection of heavily obscured sources, as it shows that they can be detected in the soft band through their soft-excess emission. This is good news for the *eRosita* instrument will conduct an all-sky imaging survey up to 10 keV. One of the primary science goals is to detect systematically all obscured accreting Black Holes in nearby galaxies and many (up to 3 million) new, distant active galactic nuclei. This technique provides a reliable tool for this purpose, especially if it is sensitive down to 0.2 keV, as this technique could be applied to sources out to $z \sim 4$.

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REFERENCES

Alexander D. M., Bauer F. E., Brandt W. N., Schneider D. P., Hornschemeier A. E., Vignali C., Barger A. J., Broos P. S., Cowie L. L., Garmire G. P., Townsley L. K., Bautz M. W., Chartas G., Sargent W. L. W., 2003, AJ, 126, 539
Antonucci R., 1993, ARA&A, 31, 473
Bassani L., Dadina M., Maiolino R., Salvati M., Risaliti G., della Ceca R., Matt G., Zamorani G., 1999, ApJS, 121, 473
Brightman M., Nandra K., 2011a, MNRAS, 413, 1206
Brightman M., Nandra K., 2011b, MNRAS, 414, 3084
Daddi E., Alexander D. M., Dickinson M., Gilli R., Renzini A., Elbaz D., CIMatti A., Chary R., Frayer D., Bauer F. E., Brandt W. N., Giavalisco M., Grogin N. A., Huynh M., Kurk J., Mignoli M., Morrison G., Pope A., Ravindranath S., 2007, ApJ, 670, 173
Della Ceca R., Castelli G., Braito V., Cagnoni I., Maccacaro T., 1999, ApJ, 524, 674
Diamond-Stanic A. M., Rieke G. H., Rigby J. R., 2009, ApJ, 698, 623
Fiore F., Grazian A., Santini P., Puccetti S., Brusa M., Feruglio C., Fontana A., Giallongo E., Comastri A., Grappioni C., Pozzi F., Zamorani G., Vignali C., 2008, ApJ, 672, 94
Gilli R., Vignali C., Mignoli M., Iwasawa K., Comastri A., Zamorani G., 2010, A&A, 519, A92
Guainazzi M., Bianchi S., 2007, MNRAS, 374, 1290
Hasinger G., 2008, A&A, 490, 905
Hopkins P. F., Hernquist L., Cox T. J., Di Matteo T., Robertson B., Springel V., 2006, ApJS, 163, 1
LaMassa S. M., Heckman T. M., Ptak A., Martins L., Wild V., Sonnentrucker P., 2010, ApJ, 720, 786
Matt G., Brandt W. N., Fabian A. C., 1996, MNRAS, 280, 823
Noguchi K., Terashima Y., Awaki H., 2009, ApJ, 705, 454
Rush B., Malkan M. A., Spinoglio L., 1993, ApJS, 89, 1
Spinoglio L., Malkan M. A., 1989, ApJ, 342, 83
Trouille L., Barger A. J., Cowie L. L., Yang Y., Mushotzky R. F., 2008, ApJS, 179, 1
Turner T. J., George I. M., Nandra K., Mushotzky R. F., 1997, ApJS, 113, 23
Ueda Y., Akiyama M., Ohta K., Miyaji T., 2003, ApJ, 598, 886
Ueda Y., Eguchi S., Terashima Y., Mushotzky R., Tueller J., Markwardt C., Gehrels N., Hashimoto Y., Potter S., 2007, ApJL, 664, L79