Title: Detecting the Gravito-magnetic field of the Dark Halo of the Milky Way (LaDaHaD)

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Abstract: We propose to locate in at least three of the Lagrange points of the Sun-Earth pair, transponders and atomic clocks aimed at exploiting the time of flight asymmetry between electromagnetic signals travelling in opposite directions along polygonal loops having the Lagrange points at their vertices. The asymmetry is due to the presence of a gravito-magnetic field partly due to the angular momentum of the Sun, partly originated by the angular momentum of the galactic dark halo in which the Milky Way is embedded. We list also various opportunities which could be associated with the main objective of the system.

Introduction

The present white paper proposes future mission(s) in the inner solar system aimed at detecting weak gravitational effects of General Relativity (GR) especially in connection with the possible existence of a galactic dark halo. The main configuration is designed for an experiment which per se requires more than one spacecraft but allows for a progressive deployment and various steps. Each step can provide essential information; furthermore, synergies are possible with other programs. For short we shall nickname our proposal LaDaHaD (Lagrange Dark Halo Detector). The main scientific objective is to look for evidence of gravito-magnetic (GM) effects due to the angular momentum of the Milky Way, both in its visible and in its dark component. As we shall see, General Relativity implies that any rotating mass distribution gives rise to a peculiar gravitational interaction resembling the effect of a magnetic field on moving charges. If the Milky Way (MW), as it is now a standard assumption, is immersed in a dark halo much more massive than the visible stars of the galaxy, it is reasonable to look for any effect originating from the rotation of such a huge halo. Considering the physical situation in the solar system, we may widen a bit the goal of the proposed mission(s) saying that the main target is any detectable gravito-magnetic effect over distances of the order of the inner solar system. Besides the mentioned angular momentum of the Milky Way, we can include the detection and measurement of the solar gravito-magnetic field. So far GM fields have been looked for at the local scale in the terrestrial environment (besides verifying the compatibility between their presence and observations of the behavior of binary systems comprised of very compact objects such as neutron stars). The mission(s) are aimed at sensitivities and accuracies better than the results obtained so far for the terrestrial effect. In general we remark that the experiment could also verify or strongly constrain deviations from GR, especially the ones expressed by the PPN parameters $\alpha_1$ and $\gamma$.

The envisaged experimental setup consists in placing spacecrafts in the Lagrange points of the Sun-Earth pair. The task assigned to each station is to bounce electromagnetic signals (pulses) along closed paths attached to the Lagrange points and to the Earth. The advantage of the Lagrange points system is that it is tied to the Earth and rotating with it; furthermore, all stations are in free fall and locally equivalent to inertial laboratories.
The essence of the experiment would be the local measurement of times of flight along the closed paths, made at one of the stations (or the Earth). More specifically the looked for data would be the time of flight differences between clock- and respectively counterclockwise beams along the same path, which are directly connected with the presence of a gravito-magnetic field and the kinematical rotation of the whole apparatus.

The envisaged setup would also allow for a zero-area Sagnac interferometer configuration at the scale of millions of kilometers, which would allow monitoring of Gravitational Waves (GW).

The same stations at the Lagrange points could act as artificial pulsars and bases of a Relativistic Positioning System usable for the guidance of space missions in the Solar System.

The main technological and methodological challenges to be tackled concern the reliable measurement of time intervals as short as $10^{-17}$ s or less, in a space environment. The whole communication and signal management system needs also to be properly designed and materially implemented. In terms of management and modeling, also the evaluation and control of the movements of the stations around the Lagrange points must be achieved.

Last but not least, sending missions to $L_4$ and $L_5$ offers the opportunity to explore the local environment of those shallow effective potential wells, by now marginally known.

The complete experimental setup could be deployed gradually, starting with the Earth and already planned (for different purposes) or even flying missions, plus one station in $L_5$. The corresponding triangle would already permit to start gathering data; then the other stations could be launched one by one, progressively enriching and upgrading the system.

Background of the proposal

General Relativity is perhaps the most elegant and important theory of contemporary physics for describing the behavior of the universe at the largest scale. Though being so successful, there are a number of major open problems: the most fundamental and embarrassing one is the deep incompatibility between GR and quantum mechanics; then we have the dark matter (DM) problem; third in the list, but not lower in importance, comes the accelerated expansion of the universe or dark energy (DE) problem. Indeed, we know that, according to the standard cosmological $\Lambda$CDM model, the dominant form of energy-mass density in the universe exists in the form of nonrelativistic (cold) dark matter and of negative pressure dark energy, while ordinary (baryonic) matter, forming all visible stars and interstellar (intergalactic) dust, contributes only less than 5% (Planck, 2018). Evidence for DM presence is known at several levels (Freese, 2017) and a substantial amount of DM is required to explain galaxies stability (Saxton, 2013), clusters of galaxies permanence, gravitational lensing (Massey, Kitching, & Richard, 2010) and spectrum of cosmic microwave background radiation. The physical nature of DM is being intensively discussed and searched for. Early investigations concentrated on ordinary matter in invisible (nonluminous) objects; nowadays, massive particles interacting only very weekly with ordinary matter (WIMPs) and/or very light axions are leading candidates. None of the recent experiments, however, has been able to detect them.

Spiral galaxies, for which rotational curves measurement could be extended into stars free edge regions via radio emission observations of neutral hydrogen (Sofue & Rubin, 2001), are the best testbeds for assessing DM scenarios. For long-term stability of such galaxies, some kind of force has to provide centripetal acceleration, which has to be precisely tallied to measured tangential (rotation) velocity at all distances from the galaxy centre. Mutual gravitational attraction of stars, central black holes and interstellar dust provide a significant part (nearly all) of the required centripetal force at small distances (up to 5 kpc) from the centre, while flat rotation curves at larger distances undeniably
point towards some other source of centripetal force. In the DM scenario this extra force is considered to be gravitational attraction of a large halo (far beyond 30 kpc in diameter) of massive nonluminous particles with isothermal spherical distribution.

The stabilizing effect of DM is currently being contemplated only as a consequence of static or slowly changing gravitational fields ($g_{00}$ component of the metric tensor). However, the presence of dynamical components ($g_{0i}$), also known as gravito-magnetic components, might successfully mimic, or at least contribute to, the effects of DM in rotating spirals. Even though the strength needed for such dynamical fields would be too large to originate only from the angular momentum of ordinary matter in galaxies. We can estimate the GM effect of the rotating visible matter in disk dominated spiral galaxies assuming that the surface brightness corresponds to ordinary (baryon) matter density in the disk. An oversimplified model treats the system as an infinitely thin disk with an exponentially decreasing mass distribution

$$\sigma(r) = \Sigma_0 \exp \left( -r / r_d \right)$$

Eq. (1)

where $r_d$ is a characteristic disk scale length and $\Sigma_0$ the surface mass density parameter. Under these conditions and for a 220 km/s peripheral velocity of the plateau of the rotation curve (which is the case of the MW), the intensity of the GM due to the visible mass would be as shown in Fig. 1 (Valko, 2019)

Fig. 1 Calculated GM field strength generated by ordinary matter of orbiting stars situated in an infinitely thin galaxy disk, with mass distribution given by Eq. (1), for a visible disk mass $4.5 \times 10^{10}$ solar masses and constant 220 km/s plateau orbital velocity. Possible central black hole and bulge contribution are not accounted for.

Calculated gravito-magnetic field strengths, shown in Fig. 1, are very weak, varying from positive values of $\sim +2.4 \times 10^{-21}$ s$^{-1}$ at bulge edge, to $\sim -1.5 \times 10^{-22}$ s$^{-1}$ at the minimum, while crossing zero, then changing sign, at approximately $2r_d$. This behaviour is clearly related to the assumed (observed) exponential ordinary mass distribution given by Eq. (1). Such feeble gravito-magnetic fields are weaker than similar fields in vicinity of typical rotating stars. For instance, the Sun generated GM field strength at Earth distance is approx. $\sim 8.3 \times 10^{-20}$ s$^{-1}$. Such fields are too weak to be sensed by any current experimental detection scheme.

The rotational velocity terms associated with the DM scenario for M31 (Carignan, Chemin, Huchtmeier, & Lockman, 2006) provide a linearly growing contribution to the rotation curve with a
slope of $\sim 1.2 \times 10^{-16} \text{ s}^{-1}$. This result could be interpreted also as the influence of a homogenous GM field, perpendicular to the plane of the disk of the galaxy, with an intensity $B_{\text{GM}} \sim 1.2 \times 10^{-16} \text{ s}^{-1}$.

Coming to our Galaxy, if the Sun should move around the centre of the Milky Way solely under the effect of a uniform GM field, ignoring for a moment purely attractive gravitational forces of both ordinary and dark matter, the field strength would have to be $B_{\text{GM}} \sim 4.5 \times 10^{-16} \text{ s}^{-1}$. On the other hand, according to a realistic MW mass distribution model (Gerhard, 2002), an additional force component is needed, to account for $30 \text{ km/s}$ of orbital velocity from the total $v_{\text{LSR}} = 220 \text{ km/s}$. The Local Standard of Rest velocity, $v_{\text{LSR}}$, is the average tangential velocity of matter in the neighbourhood of the Sun with respect to the centre of the MW. A local value of $B_{\text{GM}} \sim 1.1 \times 10^{-16} \text{ s}^{-1}$ would account for this additional centripetal force.

In fact, exploiting the analogy with classical electromagnetism, the presence of a GM field would produce on moving masses a force similar to the EM generalized Lorentz force. The induced three-acceleration $\vec{a}$ can be written as (Ruggiero & Tartaglia, 2002):

$$\vec{a} \equiv -c^2 \nabla g_{00} + 2 \vec{v} \times \vec{B}_{\text{GM}}$$

Eq. (2)

where the gradient of $g_{00}$ (time-time component of the metric tensor) is in practice the gradient of the Newtonian gravitational potential and $\vec{v}$ is the velocity of the test mass.

These fascinating enigmas and the mentioned intriguing possibilities push on fostering experimental (and theoretical) research aimed at verifying all implications of GR, both in the extremely high and in the ultralow energy domains. At the scale of the solar system, all GR effects lie in the weak and ultra-weak gravity range. The interest of science is concentrated on the confirmation of, and possible deviations from, the behavior predicted by the theory. The scope of the present proposal is to envisage the opportunity of an experimental setup and measurement technique able to evidence the presence of GM effects associated with the rotation of the dark matter halo of the Milky way (if it exists); several further information could be retrieved from such an experiment: from possible deviations from GR to the measurement of the angular momentum of the sun and more.

Experimental rationale

Among the facts worth investigation, as we have seen, there are the effects of the rotation of the source of gravity. According to GR, a mass does not simply produce attraction towards other masses, but generates also an additional interaction, dragging test bodies around with the proper rotation of the source and giving rise to a force according to Eq. (2). This contribution to the gravitational interaction is known as gravito-magnetic field or, at least in its most known form, as Lense-Thirring (LT) effect (Thirring, 1918) (Lense & Thirring, 1918). Direct measurements of LT have been so far performed in the space surrounding the Earth and having our planet as source. The effect has been verified analyzing the precession of four gyroscopes carried by a satellite in a polar orbit (GP-B experiment) with a 19% accuracy (Everitt & al., 2011). A better result (10% accuracy) has been obtained studying the orbits of the LAGEOS and LAGEOS2 satellites, reconstructed by means of laser ranging techniques, (Ciufolini & Pavlis, A Confirmation of the General Relativistic Prediction of the LenseThirring, 2004) (Ciufolini & al., Testing Gravitational Physics with Satellite Laser Ranging, 2011). The result has been improved reaching a 5% accuracy in the LARES experiment (still under way) (Ciufolini & al., A test of general relativity using the LARES and LAGEOS satellites and a GRACE Earth gravity model, 2016). After a refinement of the analysis of the usual (Newtonian) field of the Earth and of the non-gravitational perturbations of the orbit, performed by the LARASE Team, the evaluation of the accuracy has now been brought to 2% (Lucchesi, et al., 2019).

All these tests need averaging or integrating over values of the GM field of the Earth along the orbits of the probes; furthermore, in order to disentangle the proper LT effect from the ordinary gravitational
attraction, a detailed modeling of the latter around the planet is required. A different approach is envisaged by the GINGER experiment, now at the R&D stage. In this case the measurement would be local and performed at a fixed position on (or near) the surface of the Earth (Bosi & al., 2011) (Di Virgilio & al., A ring lasers array for fundamental physics, 2014) (Tartaglia, Di Virgilio, Belfi, Beverini, & Ruggiero, 2017). Besides these experiments, indirect evidence for GM effects is found in the dynamics of binary star systems and in particular in the spin-orbit coupling for the companion stars (Weisberg & Huang, 2016).

The idea we propose here is inspired to the same principle on which the ring lasers of GINGER work. If we force light to move along a path closed in space (by means of optical fibers or mirrors) and let the whole apparatus rotate, the right-handed and left-handed beams in it take different times to return to the emission point. This is the essence of the classical Sagnac effect (Sagnac, 1913); the effect, contrary to what its eponymous thought, is a relativistic one. When a gravitational field is present, besides the purely kinematic effect of rotation, additional GR terms appear in the form of effective rotations and accounting for both the coupling between rotation of the device and the ordinary gravitational attraction, and the GM field (LT drag).

Our proposal is to make an experiment in space, working at the scale of the internal solar system, but based on local time measurements: the source of the gravitational effects could be the angular momentum of the Milky Way, both of its visible and its dark component, and/or of any other possible source of GM field, including a possible primordial relic field. We expect the dark halo to rotate since the ordinary matter in the galaxy rotates. Usually the galactic dark halos are assumed to be non-rotating. The assumption has practical motivations, but would correspond to a very peculiar situation in the universe: non-rotation, if taken literally, is an absolute condition (like, locally, for spin 0 particles). However, the conjecture of co-rotation of the halo with the embedded galaxy seems reasonable.

The mutual gravitational interaction between baryonic and non-baryonic matter, considering that the galactic spirals are not homogeneous, must produce a reciprocal drag and the apparent stability of the configuration of the Milky Way makes the equal rotation rates of the two components credible. The problem is under scrutiny when referred to the rotation of gas halos of galaxies (Hodges-Kluck, Miller, & Bregman, 2016) (Tahir, De Paolis, Qadir, & Nucita, 2019) and the analogy looks reasonable. If so, the mass commonly attributed to the dark halo suggests that the associated GM effect be not entirely negligible with respect to the usual direct attraction (gravito-electric effect). Besides the mentioned possible galactic GM, another source which will in any case be present is the angular momentum of the Sun: on one side it would be interesting to measure it, in order to retrieve independent information on the internal structure of our star; on the other, a problem to solve will be how to discriminate the solar effect from the galactic contribution.

Instead of measuring a beat frequency produced by the superposition of two stationary counter-rotating electromagnetic beams, as in the case of ring lasers, the idea is now to directly measure the time of flight asymmetry between two electromagnetic pulses travelling along the same space trajectory in opposite directions. In order this idea to be implemented we need: a) a long enough path to produce a detectable time of flight difference; b) a stable, possibly non rotating device (in the reference frame of the observer) in order to remove purely kinematical effects. Now, a stable configuration is found in space when considering the Lagrange points ($L_i$) of a two body gravitationally bound system. A remarkable property of the $L$ points is that they keep (almost) fixed positions with respect to the rotating pair of masses. $L_i$’s are the five points (that is their number for any pair of objects), where the centrifugal force is balanced by the gravitational attraction towards the barycenter of the pair. Let us consider the Sun/Earth pair: a spacecraft located in one of the $L$ points will orbit the center of mass of the system (which lies within the body of the Sun) at the same angular velocity as the Earth does,
so keeping the relative position fixed. This would literally be the case only if the orbit of the Earth were perfectly circular, which is not; however, the effect of the eccentricity, as well as other perturbations and deviations are, in principle, manageable. The properties of the Lagrange points have already been, and are, exploited for science purposes, especially using \( L_2 \) (lying along the Sun-Earth line on the same side, but outwardly with respect to the Earth), then \( L_1 \) (as \( L_2 \), but internal to the Sun-Earth pair). By the way we plan to profit of the Gaia mission, positioned in \( L_2 \), (part of our team is also member of that collaboration) through the exploitation of the fully relativistic framework implemented for the data analysis of that mission in the ever-present and ever-changing solar system overlapping gravitational fields.

**The measurement strategy**

Imagine a spacecraft placed in any of the \( L \) points, and repeaters (transponders) in the others. The situation is sketched in Fig. 2.

![Fig. 2. Positions of the Lagrange points. The yellow disk is the Sun; the smaller blue disk is the Earth.](image)

From the station, you send a couple of electromagnetic pulses towards two other \( L \) points; the transponders there re-send the pulses one to the other, then back to the main station. A bit more problematic is \( L_3 \), which is behind the Sun and not visible from Earth; it would undoubtedly be more difficult to put a spacecraft there and to exchange data with it. However, \( L_3 \) is visible from \( L_4 \) and \( L_5 \), which could work as relay stations, and it could offer an interesting and advantageous configuration for the measurement we are proposing. Letting \( L_3 \) aside for a moment, a possible choice could correspond to the sequence \( L_2 - L_4 - L_5 - L_2 \) and viceversa. As said, the times of flight in opposite directions would be different. Indeed, writing the line element for a space-time having a chiral symmetry about the time axis (i.e. having a gravito-magnetic field) and with an axially symmetric central mass, one has (general coordinates)

\[
ds^2 = g_{0i} c^2 dt^2 + g_{ij} dx^i dx^j + 2c g_{0i} dx^i dt
\]

Eq. (3)

The mixed term \( g_{0i} \), in the weak field approximation leading to the GM field, may be interpreted as a space component of a potential three-vector \( \vec{h} \) similar to the vector potential of classical electromagnetism. Fully exploiting the analogy, the GM field is then:

\[
\vec{B}_{GM} = c \vec{\nabla} \times \vec{h}
\]

Eq. (4)
For electromagnetic signals it is \( ds = 0 \), so from (3) we may obtain the coordinated time of flight element \( dt \). Considering a closed path, choosing a future-ward propagation, remarking that for opposite directions propagation \( dx^i \) changes sign, taking into account the gravitational field at the emission/arrival station, the right/left proper time of flight difference for the observer, \( \delta \tau \), will be (Di Virgilio & al., A ring lasers array for fundamental physics, 2014):

\[
\delta \tau = -\frac{2}{c} \sqrt{g_{00}} \int \frac{g_{0i}}{g_{00}} \, dx^i \quad \text{Eq. (5)}
\]

In the case of a circuit including our \( L \) points, it is convenient to use heliocentric ecliptic space coordinates, where, in our case, \( \theta = 0 \) (the \( L \) points lie in the ecliptic plane). In this reference frame, considering the symmetry, it is \( g_{0i} \equiv g_{0\varphi} \), and the corresponding elementary displacement is \( dx^i \equiv rd\varphi \).

**The galactic dark halo**

Evaluating the effect of the GM field of the Milky Way, we must notice that our Lagrangian triangle (or any other Lagrangian polygon we would like to use) will be orbiting the Sun together with the Earth, hence, during the year, it will change its position with respect to the center of the Milky Way. This periodic motion implies the "device" will meet a changing galactic GM field, \( \vec{B}_{GM} \). The maximum size of the displacement is of the order of the diameter of the Earth's orbit, i.e. \( \delta l \sim 3 \times 10^{11} \, m \); on the other side, the distance of the Sun from the galactic center is roughly \( l \sim 2.5 \times 10^{20} \, m \) (Honma, 2012). It is reasonable to expect that the relative change of the intensity of \( \vec{B}_{GM} \) (seasonal variation) be (Tartaglia A., Detecting the Angular Momentum of the Galactic Dark Halo, 2019)

\[
\frac{\delta B_{GM}}{B_{GM}} \sim \frac{\delta l}{l} \sim 10^{-9} \quad \text{Eq. (6)}
\]

Actually, if we are interested in detecting the possible GM field of our galaxy with accuracies at best in the order of, say, 1% we conclude that \( B_{GM} \) is practically constant over the interior solar system. Of course \( B_{GM} \) is a vector, but we may easily extend the above conclusion to its direction also: the GM field of an axially symmetric space-time is expected to have a dipolar structure (Bosi & al., 2011), so the field is expected to be perpendicular to the galactic plane all over that plane. The Sun is thought not to exactly lie on the galactic plane, furthermore the plane of the ecliptic is at an angle \( \alpha \sim 60^\circ \) with respect to the galactic plane. Summing up, and staying within the borders of reasonable approximations, we may assume that, the whole year round, the GM field of the Milky Way is a constant vector at an angle \( \alpha \) with the ecliptic north pole.

This situation simplifies the evaluation of the possibly detectable signal in LaDaHaD. Recalling Eq. (4) we may use Stokes' theorem and write:

\[
\delta \tau_{GM} = \frac{2}{c^2} \int \vec{B}_{GM} \cdot \hat{n} \, dS
\]

The flux integral is over the area, \( S_L \), contoured by the Lagrangian triangle (or other polygon). The practical local uniformity of \( \vec{B}_{GM} \) immediately gives the result of the integral:

\[
\delta \tau_{GM} \approx \frac{2}{c^2} B_{GM} S_L \cos \alpha \quad \text{Eq. (7)}
\]

We may reverse Eq. (7) to get \( B_{GM} \):
The area of the $L_4 - L_5 - L_2$ triangle is $\sim 9.9 \times 10^{21}$ m$^2$. Let us suppose the smallest detectable time of flight asymmetry be $\sim 10^{-16}$ s. In such conditions we should be able to sense GM fields of the Milky Way at the position of the solar system

$$B_{GM} \geq 3 \times 10^{-22} \text{ s}^{-1}$$

Of course, what we would measure with LaDaHaD would indeed combine the effect of the visible matter and of the dark halo. Just to give an idea of the relative importance of the two contributions let us recall that GM fields originating from rotations are in general proportional to the angular momentum of the source. A rough guess can then be obtained comparing the angular momentum of a homogeneous disk (symbolizing baryonic matter) with that of a homogeneous sphere (the dark halo) rigidly rotating at the same angular speed. The ratio of the latter to the former is just, modulo a factor of order 1, the ratio between the two masses, which, in the case of the Milky Way could be in the order of 10. In other words, the contribution of the dark halo would clearly prevail on the one of ordinary matter (which has been shown in Fig. 1).

**Solar gravitomagnetism**

Considering the only GM field of the Sun (neglecting the small perturbation from the Earth) and recalling Eq. (3) for an axially symmetric line element in vacuo and weak field approximation, it is:

$$g_{0\phi} = \frac{2GJ}{c^3 r^2} \quad \text{Eq. (8)}$$

$J$ is the angular momentum of the Sun.

In the reference frame of the observer (main station) rotating with the Earth and with the whole set of $L$ points, introducing the astronomical unit $\rho$ (Sun-Earth average distance $\sim 1.5 \times 10^{11}$ m), assuming weak field conditions and keeping terms up to the lowest order in the angular momentum, Eq. (8) becomes:

$$g_{0\phi,\text{main}} = -\frac{GM}{c^2 \rho} \left[ \frac{\rho}{r} + \left( 1 - \frac{\rho}{r} \right) \frac{GM}{c^2 r} \right] + 2 \frac{GJ}{c^3 r^2} \quad \text{Eq. (9)}$$

$M$ is of course the mass of the Sun.

Let us now introduce numbers. The distance of $L_4$ and $L_5$ from the Earth (one ahead and the other behind, along the terrestrial orbit) is $\rho$. The distance of $L_1$ and $L_2$ (inward and outward along the Sun-Earth line) is $\sim 1.5 \times 10^9$ m. Using Eq. (5), with $g_{00} = \sqrt{1 - \frac{2GM}{c^2 \rho}}$, in the case of an $L_2 - L_4 - L_5 - L_2$ circuit, we see that the total time of flight turns out to be $\sim 2000$ seconds and the right/left asymmetry is (Tartaglia, et al., 2018)

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1 The triangle is treated as being Euclidean and the curvature of light rays due to the solar lensing is much smaller than the effects we are interested in.
If one succeeded in including $L_3$, the triangle $L_3 - L_4 - L_5$ would correspond to a time of flight in the order of $8000 \text{ s}$ and it would be $\delta \tau \gtrsim 10^{-12} \text{ s}$.

By the way, if the galactic GM field able to account for the excess $v_{LSR}$ would actually exist, the time of flight difference would be of the order of $\sim 10^{-11} \text{ s}$ ($\sim 10^{-10} \text{ s}$, if $L_3$ were used), already accounting for the angle between the ecliptic and the MW disk planes.

If the spacecraft located in $L_2$ (or any other point) is equipped with an atomic clock activated by the arrival of the first pulse and stopped by the arrival of the second, the above time difference may be measured at the 1% level (or better). In this way a test of the GR effects in the field of the Sun is performed, possibly with accuracies better than those attainable in the terrestrial environment.

The advantage of using the system of the Lagrange points, in addition to the size of the loop, as said, is that it constitutes a stable frame co-rotating with the Earth, so that the experimental setup is also stable and the measurement may be repeated many many times in the same geometrical conditions, allowing to average out random disturbances. Furthermore, the objects lying in the $L$’s are in free fall, and so is the whole setup, allowing to get rid of kinematical, non-inertial, disturbing or even dominant, effects, typical of terrestrial laboratories.

It is the case to remark that solar LT field measurements, especially its dynamics, could be an interesting tool to address the Sun interior rotation profile, particularly within the convective zone. Once more the most effective configuration would include $L_3$.

**Synergy with LISA**

The measurement strategy proposed for LaDaHaD offers an interesting opportunity of synergy with LISA. As known, LISA is a space interferometer designed to reveal gravitational waves (GW). After the recent detection of a number of signals originated by the merger of pairs of black holes (Abbott & al., 2016) or even pairs of neutron stars (Abbott & al., 2017), the international effort to further detections has become even more intense than before.

The configuration of LISA is an equilateral triangle $2.5 \times 10^9 \text{ m}$ in side having freely falling transponders at the end of each side (Danzmann, 2017). The center of the triangle would orbit the sun along the orbit of the Earth at some 50–65 million km behind our planet. Laser beams along the arms of LISA would serve both as GW sensors and as controllers of the geometry. It is evident that the closed contour could also be used for Sagnac-like experiments like for LaDaHaD: the LISA triangle would have an area four orders of magnitude smaller than the LaDaHaD triangle but at the same time it would lead to a seasonal modulation of the inclination angle of the plane of the polygon with respect to the axis of the Milky Way. A collaboration with LISA would be both a precious intermediate step towards the full deployment of LaDaHaD and a relevant autonomous investigation of the GM field of the MW and of the Sun.

**Gravitational Waves**

Besides the mentioned possible synergy with LISA, LaDaHaD could also lend an interesting complementary opportunity for the detection of GW. In fact, there is an approach that can employ the same setup and method described for the GM measurement. Recently there has been a proposal to exploit (in a circum-terrestrial environment) a Sagnac-like measurement along a triangular loop in
space, for detecting GW originated from intermediate mass black hole mergers (Lacour & al., 2018). Already in the ‘90s of the past century Sun, Fejer, Gustafson and Byer (Sun, Fejer, Gustafson, & Byer, 1996), starting from a previous idea of Weiss, proposed and analyzed a so called zero-area Sagnac interferometer for the detection of GW. The zero-area Sagnac interferometer organizes the path of light beams so that the contoured area is zero. In practice, one has a pair of twisted half loops and each beam travels along the two loops once clockwise, once counter-clockwise. The consequence is that the signal of the kinematic and of the effective rotations cancel; if, however, the strain induced by a GW is present, the transient difference in the length of the arms of the interferometer maintains its effect and is recorded as a residual asymmetry in the phases of the beams. On Earth the principle works, but the storage time in one arm of the interferometer can hardly be longer than a ms. In the case of our Lagrangian system a zero-area configuration could be achieved including earth-based stations and using for instance an $L_4 - L_2 - Earth - L_5 - L_2 - Earth - L_4$ circuit. In that case the “storage time” (here the flight time) in a half loop would be in the order of one thousand seconds. The angle between the “arms” (important in order the strain of the GW to be manifested) would be $120^\circ$ rather than $90^\circ$, but the permanence time would be $10^6$ times longer than in a terrestrial device.

**Relativistic positioning**

Together with the direct scientific return provided by LaDaHaD in evidencing GR effects at the scale of the solar system, the whole set and the same principle could provide an important support to any other mission in space outside the Earth environment. One of the delicate problems concerning all such missions is the localization and guidance of the spacecraft. Such task is currently performed by a combination of different techniques, including self-guidance through the image of the sky. Most often, however, and for most of the time, the mission control is carried out from Earth. There, ranging methods give usually good results, however the transverse positioning progressively worsens while the distance increases. The situation could be much better with beacons located in the Sun-Earth Lagrange points, like in the configuration proposed for testing the GR effects. In fact, not much more gears would be needed than what already required by the main experiment: each spacecraft of the LaDaHaD system should act as a stable emitter either of a continuous electromagnetic wave at a fixed frequency or of pulses issued at a fixed rate. The emission should be isotropic or at least over wide angles. The positioning method based on local time measurements autonomously performed by the user (the object to be positioned) is intrinsically relativistic (we may call it Relativistic Positioning System: RPS) and has been proposed and discussed in various papers, and in particular in (Tartaglia, Ruggiero, & Capolongo, A null frame for spacetime positioning by means of pulsating sources, 2011).
In short, the idea is as follows. We start from a physical basis composed of (at least) four independent emitters of electromagnetic pulses (or continuous harmonic waves) whose emission rate is as stable as possible, and is known. The emitters form a physical frame with respect to which the positioning is made. In four-dimensions the wave-fronts of the pulses from the sources cover space-time with a regular array made of “cages” whose edges are parallel to the propagation directions from the emitters and whose size (when splitting space and time) is the proper wavelength (or distance from successive pulses) and the proper emission period. In such four-dimensional grid, a traveler is represented by a continuous line (its world-line) crossing the boundaries of successive “cages” (see Fig. 3, drawn in two, i.e. 1+1, dimensions); the situation is described and visually depicted in (Tartaglia A., Relativistic space-time positioning: principles and strategies, 2013). While receiving the signals from the various sources, the traveler must have the possibility to distinguish and identify each source, then he/she counts the pulses and measures on his own clock the time intervals between the arrivals of one pulse and the next. If the signals are continuous, equivalently, he/she takes note of the relative phase between the incoming waves. The measured time spans are proportional to the intervals (in the GR sense) between successive arrival events. Simple geometrical considerations allow the user to reconstruct, from the measured sequences of arrival times, its own coordinates along the light cones, then, more familiarly, in the reference frame of the emitters. Such RPS has been thought and tested using pulsars as beacons (Ruggiero, Capolongo, & Tartaglia, 2011). In the case of LaDaHaD, the emitters in the $L$ points would act as “artificial pulsars”. Once more, the advantages of the Lagrange points is that they form a reference frame co-moving with the Earth. Furthermore, the size of the base is such that geometric dilution effects can be minimized.

**Relativistic modelling of the solar system**

The advancement in astronomical observations and instrumentation requires reconstructing light propagation at a very high level of precision, whenever the accuracy of the measurements, in terms of lengths, is comparable to the square root of the inverse curvature of the background geometry due to the distribution of the sources of gravity. This is particularly needed for space missions such as LaDaHaD and Gaia, which necessarily require extremely accurate astrometric observations modeled within a fully, suitably accurate, relativistic framework.

The new model for ‘tracing’ photons developed for the Gaia mission (see (Crosta, Tracing light propagation to the intrinsic accuracy of spacetime geometry 28, 2011), (Crosta & al., The ray tracing
analytical solution within the RAMOD framework. The case of a Gaia-like observer (2015) and references therein; (Vecchiato & al., 2012)) is now part of the European excellence in the field of modeling photon paths in low gravity and of the Gaia effort of building, with two independent methodologies the first fully relativistic reconstruction, both differential and absolute, of the celestial sphere. The other approach, GREM, is that in (Klioner, A Practical Relativistic Model for Microarcsecond Astrometry in Space, 2003).

In particular, RAMOD (“Relativistic Astrometric MODel”, (Bini, Crosta, & de Felice, 2003), (de Felice & al., A General Relativistic Model of Light Propagation in the Gravitational Field of the Solar System: The Static Case, 2004) (de Felice & al., 2006), (Crosta & Vecchiato, Gaia relativistic astrometric models. I. Proper stellar direction and aberration, 2010)) naturally recovers the seminal and fundamental works of Kopeikin & Schäfer (Kopeikin & Schäfer, 1999), and Klioner ( (Klioner, A Practical Relativistic Model for Microarcsecond Astrometry in Space, 2003) and references therein) and provides a fully general-relativistic analysis of the inverse ray-tracing problem based upon a measurement protocol in GR (de Felice & Bini, Classical Measurements in Curved Space-Times, 2010). The latter uses a 3+1 characterization of space-time in order to measure physical phenomena along the proper time and on the rest-space of a set of fiducial observers. Then, the unknown of RAMOD is the four-dimensional local line-of-sight as measured by the fiducial observer at the point of observation. While the physical spatial component is used for the Gaia data modeling and compared with the coordinated one used in GREM, the time component can be exploited for the LaDaHaD measurements. Moreover, the RAMOD method is close (Bertone & al., 2014) to the Time Transfer Function formalism (TTF, (Teyssandier & Le Poncin-Lafitte, 2008), (Le Poncin-Lafitte & Teyssandier, 2008)) that stands as a development of Synge’s World Function, an integral approach based on the principle of minimal action. While the World Function is an implicit equation of the photon trajectory nearly impossible to solve, the TTF formalism gives up some generality to provide important information about the propagation of an electromagnetic wave between two points at finite distance from one another.

This scientific activity has not only contributed to establish novel theoretical and numerical methods for treating light propagation consistently with the covariant approach of GR, but has also provided new strategies to go further in testing gravity and its connection with the scientific potential of a local measurement entangled in the varying gravitational fields of the Solar System. Any inconsistency in the relativistic model(s) would invalidate the quality and reliability of the estimates, indeed all the relevant scientific outputs, and, in addition, it can be the clue for uncovering new effects. In this regard, LaDaHaD can offer the unique opportunity to complete the theoretical comparison of the existing approaches and set the scientific case for further developments and applications, even in the quantum gravity regime (Khavkine, 2012).

More opportunities

Pulsar timing

The use of the emitters as “artificial pulsars” located at well-defined positions in space lends an opportunity also for astrophysicists. In fact, the manmade “pulsars” may be used to help timing the real ones.

Gravitational potential mapping at the L points

Studying the behavior of the spacecraft in each L point can be a method to map the corresponding gravitational potential. The movements of a station around its own L point can indeed be reconstructed using the other stations as a reference frame and applying the RPS. The reconstructed details of the movement give a map of the local effective potential in the vicinity of the L point.
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Recognition of the space environment of $L_4$ and $L_5$

$L_4$ and $L_5$ are the most distant Lagrange points (not considering $L_3$), so until now not much attention has been payed to their environment, even though it would probably be quite interesting. ESA is indeed planning to send a mission to $L_5$ that could be combined with LaDaHaD. $L_4$ and $L_5$ are (weakly) stable, being local (shallow) minima of the effective potential (including Coriolis forces). This feature is favorable for positioning there the transponders and stations of LaDaHaD; in the same time, the local potential well may imply the presence of celestial bodies trapped there. It is already known that at least one asteroid (2010 TK$_7$) is dwelling in $L_4$ (Connors, Wieger, & Veillet, 2011): it is the first “Trojan asteroid” of the Earth to be discovered. More could be there and in $L_5$ too. Missions sent to lay transponders/beacons of LaDaHaD in $L_4$ and $L_5$ would offer the possibility to explore the space around those points.

Scientific and technological requirements

In order to verify the practical possibility to carry the present proposal to completion, a number of detailed studies both on the scientific and on the technological aspects of the related missions are necessary.

On the side of scientific and conceptual tools a list of requirements follows.

a) The solar system is of course not made of the Sun-Earth pair only. An analysis of the gravitational perturbations induced by the other celestial bodies is in order, both for the stability of the positions of the Lagrange points and for the influence on the signal that would be measured.

b) All effects have to be analyzed from the view point of GR in order to build a consistent interpretation frame for the data in view of the expected results.

c) A detailed study of the effects of deviations from GR in the framework of PPN formalism is in order. An analysis is already present in (Bosi & al., 2011) and the relevant parameters taken into account there have been $\alpha$ (accounting for the presence and relevance of preferred reference frames) and $\gamma$ (accounting for the effect of spatial curvature alone). The possible relevance of such and other parameters must be discussed at the scale of the inner solar system. Furthermore, the possible presence of galactic or even relic GM fields must be considered and modeled in order to provide appropriate templates for the interpretation of the results.

d) Being the positioning at the $L$ points crucial, it is of paramount importance the study of the orbits of the spacecraft of the experiment around those points and the related dynamics and stability. Libration movements are expected in $L_4$ and $L_5$, where the orbits can be stable, whereas weakly unstable Lissajous orbits are typical of $L_1$ and $L_2$ (saddle points of the effective potential). Stabilization strategies must be envisaged. One should also account for the proper motion of the $L$’s with respect to the Earth, due to the eccentricity of the orbit of our planet.

e) A model must be worked out accounting for the effects of the movements of the stations, according to point d), on the final evaluation of the time of flight difference along closed paths among the $L$ points. Semi-analytical investigations of the main periodic orbits around the Lagrange points (Celletti, Pucacco, & Stella, 2015) (Ceccaroni, Celletti, & Pucacco, 2016) (Páez, Locatelli, & Efthymiopoulos, 2016) provide accurate approximations to model distance variations. They can also be used as seeds for numerical surveys of the unstable dynamics around $L_1$ and $L_2$.
Coming to technological needs, many aspects require in depth analyses and improvements with regard to their present implementations.

a) The core measurement of LaDaHaD concerns time. As stated in previous sections, the interval between the arrivals of the clock- and respectively counterclockwise pulses is in the order of $10^{-13}$ s (in the best case $10^{-12}$ s). In order to attain a 1% accuracy, the clock (and the measurement process) must have a sensitivity of at least $10^{-15}$ s. This is not a problem in terrestrial laboratories, where one can achieve results one or two orders of magnitude better than that. One has, however, to think of atomic clocks in space, onboard an observatory located in one of the $L$ points (such as $L_1$ or $L_2$), endowed with the required accuracy and a stability sufficient for a long mission (not less than months). A possible solution in order to ensure the best accuracy and for the longest time would be to permanently and for all configurations (see ahead the deployment strategy) keep an Earth laboratory as hub of our experimental polygonal network. In any case, the need for having and controlling clocks and oscillators in space cannot be dispensed of completely, so that the problem must be carefully studied.

b) The main station must be equipped also with an appropriate elaboration capacity. Even though the requirement there is not much stringent: the algorithm to be used, in the case of RPS, is linear. More important is the onboard data storage capacity, connected with an effective communication channel for transferring information to the Earth.

c) The technology of the transponders to be used at the other (with respect to the main one) stations has to be analyzed carefully. Each transponder is expected to work as much as possible like a mirror, with a fixed and as short as possible delay between the arrival and the bounce back of the incoming pulses. The same problem has to be tackled for the phase-locked transponders of LISA.

d) In case of the use of LaDaHaD as a base for positioning and navigation in the solar system, each $L$ point must host an emitter of regular pulses, possibly combined with the transponder. Each emitter needs be controlled by an atomic clock; the main constraints are not on the frequency (GHz or even MHz repetition times for the pulses are sufficient), but rather on the stability over months or even years, if the RPS has to be considered as a permanent equipment for navigation.

e) Special attention must be paid to the design of all antennas installed onboard.

f) All electronic equipment, at all stations, of course needs an appropriate power supply, whose main components are solar cells and storage capacity.

**Deployment and measurement strategy**

As already explained, the core of the experiment is the local time measurement of time of flight differences along closed paths in space, based on the Sun/Earth $L$ points. The full deployment of the system is of course burdensome and requiring a long time to be completed. It is however possible to define a strategy relying on progressive steps, allowing data collection from the beginning. As previously mentioned, the $L_2$ and $L_1$ points already host various space missions. In particular, in $L_2$ we find Gaia, which is equipped for communicating with the Earth and could technically perform tasks much like the ones demanded to the proper stations of LaDaHaD. The planned duration of the mission is close to the end, but, most likely, the spacecraft can live longer than the official duration of the mission, especially if the assignments are simplified and reduced to act as a bounce back station for LaDaHaD. Furthermore, the Earth as such may be included as a vertex of the space polygons we want to use; of course in case of ground based stations one has to include the effect of the diurnal rotation
of the planet, which implies a discontinuous operation and the need to account for a peculiar relative motion.

Relying on the cited pair of bases, a first step (not a trivial one, in any case) could consist in sending a mission to $L_5$ (or $L_4$), then the triangle Earth – $L_2$ – $L_5$ would be available for data taking. An advantage of having a laboratory on the ground would be the possibility to use the best atomic clocks, whose accuracy would in principle permit the detection of solar GM terms at levels better than 1 part in $10^3$. Of course, the perturbations induced by the atmosphere should be carefully analyzed, together with the presence, close to the ground station, of the terrestrial GM field.

The next step could be to place one additional spacecraft in $L_1$. A new triangle would then be at hands, giving a useful degree of redundancy. The whole system would contain two possible intertwined loops, in the shape of the zero-area Sagnac interferometer mentioned in the GW section. That particular configuration has also been called “butterfly wings”; in this case the “wings” would be in almost closed position and the time of flight along each loop (each “wing”) would be $\sim 600$ s, with an angle between the long and the short arm equal to $60^\circ$. An improvement of this configuration would be achieved replacing Gaia in $L_2$ with a new dedicated station.

One further improvement would be realized sending a station in $L_4$. Now the possible loops would be seven (including the Earth) offering redundancy and alternative strategies. For GW monitoring, the wings of the butterfly would now be open and symmetric with respect to the Sun-Earth axis rather than the Earth – $L_5$ direction.

The complete network of stations would be attained sending a spacecraft also in $L_3$. This would represent the most sensitive configuration for the measurement of the solar gravito-magnetism and for the study of the internal dynamics of the Sun. The highest sensitivity for deviations from GR would also be attained.
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Call for White Papers for the Voyage 2050 long-term plan in the ESA Science Programme

LaDaHaD

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