Meromorphic continuation of Selberg zeta functions with twists having non-expanding cusp monodromy

Ksenia Fedosova · Anke Pohl

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Abstract
We initiate the study of Selberg zeta functions $Z_{\Gamma, \chi}$ for geometrically finite Fuchsian groups $\Gamma$ and finite-dimensional representations $\chi$ with non-expanding cusp monodromy. We show that for all choices of $(\Gamma, \chi)$, the Selberg zeta function $Z_{\Gamma, \chi}$ converges on some half-plane in $\mathbb{C}$. In addition, under the assumption that $\Gamma$ admits a strict transfer operator approach, we show that $Z_{\Gamma, \chi}$ extends meromorphically to all of $\mathbb{C}$.

Keywords Selberg zeta function · Meromorphic continuation · Non-unitary representation · Non-expanding cusp monodromy · Thermodynamic formalism · Transfer operator

Mathematics Subject Classification Primary 11M36; Secondary 37C30 · 37D35

1 Introduction
Let $\Gamma$ be a geometrically finite Fuchsian group, let $\chi : \Gamma \rightarrow U(V)$ be a unitary representation of $\Gamma$ on a finite-dimensional unitary space $V$, and let $\mathbb{H}$ denote the hyperbolic plane. For $s \in \mathbb{C}$, $\text{Re } s \gg 1$, the (twisted) Selberg zeta function $Z_{\Gamma, \chi}$ for $(\Gamma, \chi)$ takes the form

Ksenia Fedosova
ksenia.fedosova@math.uni-freiburg.de

1 Mathematisches Institut, Albert-Ludwigs-Universität Freiburg, Eckerstr. 1, 79104 Freiburg im Breisgau, Germany

Anke Pohl
apohl@uni-bremen.de

2 Department 3 – Mathematics, University of Bremen, Bibliothekstr. 5, 28359 Bremen, Germany
\[ Z_{\Gamma, \chi}(s) = \prod_{\hat{\gamma}} \prod_{k=0}^{\infty} \det(1 - \chi(g_{\hat{\gamma}})e^{-(s+k)\ell(\hat{\gamma})}) , \]  

(1)

where 1 denotes the identity map on \( V \), \( \hat{\gamma} \) in the first product ranges over all primitive periodic geodesics of \( \Gamma \setminus \mathbb{H} \), the length of \( \hat{\gamma} \) is denoted by \( \ell(\hat{\gamma}) \), and \( g_{\hat{\gamma}} \) is a primitive hyperbolic element in \( \Gamma \) that is associated to \( \hat{\gamma} \). We refer to Sect. 2 below for more details and an alternative, more algebraic version of (1).

For various combinations \((\Gamma, \chi)\) the infinite product (1) is known to admit a meromorphic continuation to all of \( \mathbb{C} \), see [21,53,57] and the more extensive references further below.

Selberg zeta functions (i.e., the infinite products (1) and their meromorphic continuations) have many applications in various areas of mathematics, in particular in spectral theory, harmonic analysis, number theory, and in mathematical physics. The arguably most important property is that the zeros of the Selberg zeta function \( Z_{\Gamma, \chi} \) are closely related to the eigenvalues and resonances of the Laplace–Beltrami operator \( \Delta \) as acting on spaces of \((\Gamma, \chi)\)-automorphic functions on \( \mathbb{H} \) of certain regularity. This relation allows for rich results on the (non-)existence and distribution of the spectral parameters of \( \Delta \), and it establishes a link between certain geometric properties (e.g., geodesic length spectrum, volume, number of cusps; ‘classical mechanical’ aspects) and some spectral properties (e.g., eigenvalues, scattering resonances; ‘quantum’ aspects) of \( \Gamma \setminus \mathbb{H} \).

It is natural to ask to which extent these results can be generalized if in (1) the representation \( \chi \) is allowed to be non-unitary. First investigations in this direction have been conducted by Müller [44], Monheim [42], Deitmar–Monheim [8,9], Deitmar [7], Spilioti [55] and Fedosova [14].

Müller [44] considers compact Riemannian locally symmetric spaces \( \Gamma \setminus G/K \) (here, \( G \) is a connected real semisimple Lie group of non-compact type with finite center, \( K \) is a maximal compact subgroup of \( G \), and \( \Gamma \) is a \textit{torsion-free cocompact} lattice in \( G \)) and arbitrary finite-dimensional representations \( \chi : \Gamma \to \text{GL}(V) \). In addition, he allows unitary twists of \( K \), which we do not discuss here. He establishes an analogue of the Selberg \textit{trace} formula (which, for unitary representations, is known to be closely related to the Selberg zeta function) where the flat twisted connection Laplace operator \( \Delta^\#_\chi \) takes the role of the Laplacian \( \Delta \). Under the same conditions on \( G, \Gamma \) and \( \chi \), Deitmar and Monheim [8] (see also Monheim’s dissertation thesis [42]) provide a non-unitary Selberg trace formula for \( \Gamma \setminus G \). In [7] Deitmar generalizes these results to arbitrary locally compact groups \( G \).

In [9,42] Deitmar and Monheim consider hyperbolic surfaces \( \Gamma \setminus \mathbb{H} \) of \textit{finite area} (thus, \( G = \text{PSL}_2(\mathbb{R}) \), \( K = \text{PSO}(2) \) and \( \Gamma \) is a \textit{cofinite}, not necessarily cocompact

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1 To our knowledge, the cases treated in the literature are

- \( \chi \) is the trivial one-dimensional representation and \( \Gamma \) is any geometrically finite Fuchsian group,
- \( \Gamma \) is cofinite and \( \chi \) is any unitary representation,
- some combinations of non-cofinit \( \Gamma \) (e.g., Fuchsian Schottky groups) and non-trivial unitary representations \( \chi \).

Even though some combinations \((\Gamma, \chi)\) with \( \chi \) being unitary seem to be omitted in the literature, we suppose that these omissions can easily be handled with the existing methods.
Fuchsian group) and finite-dimensional representations $\chi : \Gamma \to \text{GL}(V)$ that are unitary in the cusps, i.e., for each parabolic element $p \in \Gamma$, the endomorphism $\chi(p)$ is unitary. They study convergence and meromorphic continuability for $\chi$-twisted Eisenstein series.

Spilioti [55] considers compact hyperbolic manifolds $\Gamma \backslash \mathbb{H}^n$ of odd dimension $n$ (thus, $G = \text{SO}(n, 1)$, $K = \text{SO}(n)$, $n$ odd, and $\Gamma$ is a torsion-free cocompact lattice in $G$) and arbitrary finite-dimensional representations $\chi : \Gamma \to \text{GL}(V)$. Similar to Müller, she allows an additional unitary twist of the central elements in $K$. She shows convergence of the twisted (Ruelle and) Selberg zeta functions on certain half-planes in $\mathbb{C}$. Taking advantage of Müller’s twisted Selberg trace formula she then proves meromorphic continuability of the twisted Selberg zeta functions to all of $\mathbb{C}$ and provides a spectral interpretation of their singularities. Fedosova [14] generalizes these results to arbitrary (i.e., not necessarily torsion-free) cocompact lattices $\Gamma$ in $G$.

In this paper we set out to study Selberg zeta functions $Z_{\Gamma,\chi}$ for arbitrary geometrically finite Fuchsian groups $\Gamma$ (including the non-cofinite ones) and finite-dimensional representations $\chi : \Gamma \to \text{GL}(V)$ which satisfy that for each parabolic element $p \in \Gamma$ each eigenvalue of the endomorphism $\chi(p)$ has absolute value 1. This property of $\chi$ was coined ‘non-expanding cusp monodromy’ by Eskin et al. [13] who prove lower bounds for Lyapunov exponents of flat bundles on curves, where the flat bundles are determined by such representations. Representations with this property also appeared in investigations of vector-valued and generalized modular forms [25–29,31–34].

For cocompact lattices $\Gamma$ the class of representations with non-expanding cusp monodromy is identical to that of all finite-dimensional representations of $\Gamma$. For non-cocompact Fuchsian groups this class subsumes all unitary representations, all representations that are unitary at cusps, all representations that are unipotent on all parabolic elements of $\Gamma$, as well as the restrictions to $\Gamma$ of the finite-dimensional representations of $\text{SL}_2(\mathbb{R})$, and it contains additional representations. We refer to Sect. 5 below for a more detailed discussion.

In this generality (in particular with regard to the class of representations), the convergence of $Z_{\Gamma,\chi}$ has not been studied before, and the classical proof methods in their current state of art do not apply. Therefore, as our first main result we will establish its convergence on some right half-plane of $\mathbb{C}$ (see Theorem 3.1 below).

**Theorem A** Let $\Gamma$ be a geometrically finite Fuchsian group and $\chi : \Gamma \to \text{GL}(V)$ a finite-dimensional representation with non-expanding cusp monodromy. Then the infinite product (1) converges for $\text{Re } s \gg 1$.

In addition we will prove that Theorem A is sharp as stated, it cannot be generalized to a larger class of finite-dimensional representations.

**Proposition B** (Proposition 6.1 below) Let $\Gamma$ be a geometrically finite Fuchsian group, and let $\chi : \Gamma \to \text{GL}(V)$ be a finite-dimensional representation of $\Gamma$ that does not have non-expanding cusp monodromy. Then the infinite product (1) does not converge absolutely for any $s \in \mathbb{C}$.

The combination of Theorem A and Proposition B shows that the ultimate realm of finite-dimensional twists for Selberg zeta functions for Fuchsian groups (without
using any additional type of regularization) is that of finite-dimensional representations with non-expanding cusp monodromy. Beyond this class, not even the first step in the definition of a twisted Selberg zeta function is tenable.

Other interesting directions of research are those of Selberg zeta functions with twists by infinite-dimensional representations, or for non-Fuchsian groups, or the extendability of the realm of twists by additional regularizations. These, however, will not be investigated here.

A crucial part—of independent interest—of the proof of Theorem A is to show that for every hyperbolic element \( h \in \Gamma \), the trace of \( \chi(h) \) grows at most exponentially with the length \( \ell(h) \) of the periodic geodesic on \( \Gamma \backslash \mathbb{H} \) associated to \( h \) (or, equivalently, at most polynomially with the norm \( N(h) \) of \( h \)).

**Proposition C** (Proposition 3.2 below) There exists \( c \geq 0 \) such that uniformly for all hyperbolic elements \( h \in \Gamma \),

\[
| \text{tr} \chi(h) | \ll e^{c \ell(h)},
\]

and, even stronger,

\[
\inf_{g \in \Gamma} \| \chi(ghg^{-1}) \| \ll e^{c \ell(h)}.
\]

For cocompact \( \Gamma \), Proposition C follows from the Lipschitz-equivalence of the metric on \( \Gamma \) induced by the Riemannian metric on \( \text{PSL}_2(\mathbb{R}) \) and the word metric of \( \Gamma \) (induced by any finite set of generators of \( \Gamma \)), see [36,55]. If \( \Gamma \) is cofinite and \( \chi \) is uniform at the cusps, then Proposition C has essentially already been proven in [9] (using [9, Propositions 2.7 and 2.9] in place of (13) below). For the general case of arbitrary geometrically finite Fuchsian groups \( \Gamma \) and finite-dimensional representations \( \chi \) with non-expanding cusp monodromy, the proofs in [9,36,55] do not apply anymore. For this generality, it is essential to understand which weight \( \chi \) attributes to periodic geodesics which travel ‘high into the cusps’ (long cusp excursions) or, equivalently, to understand the contribution of parabolic elements to the growth of \( \chi \). In Sect. 3 below we will provide a detailed study, including a kind of logarithm law (see Lemma 3.3 below).

Proposition C allows us to bound the twisted Selberg zeta function \( Z_{\Gamma,\chi} \) by a shift of the non-twisted Selberg zeta function \( Z_{\Gamma,1} \) for the trivial one-dimensional representation 1. The well-known convergence properties of \( Z_{\Gamma,1} \) then yield Theorem A.

The proof of Proposition C, and hence of Theorem A, is uniform for all pairs \( (\Gamma, \chi) \) of geometrically finite Fuchsian groups \( \Gamma \) and finite-dimensional representations \( \chi \) with non-expanding cusp monodromy.

Our second main result concerns the meromorphic continuability of the infinite product (1). As already mentioned above, prior to this paper, meromorphic continuability of \( Z_{\Gamma,\chi} \) has been known for unitary representations only (note that [14,55] consider odd-dimensional spaces only, thus they do not treat Fuchsian groups). For these investigations several methods (e.g., trace formulas, microlocal analysis, transfer operator techniques) have been employed and applied to various classes of \( (\Gamma, \chi) \),
resulting in alternative or complementary proofs of meromorphic extendability in different generality. In particular the usage of trace formulas and microlocal analysis requires to restrict the uniformity of the considerations to certain classes of Fuchsian groups, e.g., cocompact or cofinite Fuchsian groups or Fuchsian Schottky groups. We refer to [5,10,11,16,18,21,23,38,39,45,52,53,56,57] for examples and more details.

The proof of meromorphic continuability for Selberg zeta functions with non-unitary twists in [14,55] uses harmonic analysis on symmetric spaces, in particular trace formulas and orbital integrals. Due to this approach their results are currently restricted to (compact) hyperbolic spaces/orbifolds of odd dimension, a major problem being to guarantee that residues are integral. It would certainly be interesting to see if these methods can be adapted to even-dimensional spaces.

Also the results in [8,42,44] on the existence of Selberg trace formulas twisted by non-unitary representations (for compact spaces) or by representations that are unitary at the cusps (for non-compact hyperbolic surfaces) rely on harmonic analysis of symmetric spaces. It would certainly be interesting to understand if these trace formulas integrate to Selberg zeta functions and to which extent they generalize to non-compact spaces, spaces of non-finite volume as well as beyond representations that are unitary at the cusps.

Nevertheless, in order to show the meromorphic continuability of (1), we here will develop transfer operator techniques which permit a uniform approach for all admissible combinations \((\Gamma, \chi)\). In particular, in contrast to some other methods, we will not need to distinguish between Fuchsian groups with and without cusps, or cocompact, cofinite and non-cofinite ones.

However, we will suppose that \(\Gamma\) admits a strict transfer operator approach. This means, roughly, that there exists a transfer operator family \(L_{\Gamma,s}\) such that

\[
Z_{\Gamma,1}(s) = \det \left( 1 - L_{\Gamma,s} \right), \quad \text{Re} \, s \gg 1.
\]

In other words, the Selberg zeta function \(Z_{\Gamma,1}\) for the trivial one-dimensional representation \(1\) is represented by the Fredholm determinant of a transfer operator family on some right half-plane in \(\mathbb{C}\) where \(Z_{\Gamma,1}\) is given by the infinite product (1). We refer to Sect. 4.2 below for a precise definition.

In view of the already existing strict transfer operator approaches and the different methods for their construction [1,3,4,6,16,17,22,37–39,41,43,46,47,50–52,54], it might well be that this requirement is not a severe restriction on \(\Gamma\) at all. Moreover, it is most likely that with the methods we propose in this article the meromorphic continuability of (1) can also be shown starting with non-strict transfer operator approaches as provided in [40,43,51] (for certain classes of cofinite Fuchsian groups). In the latter case, the representation of \(Z_{\Gamma,1}\) is of the form

\[
Z_{\Gamma,1}(s) = h(s) \det \left( 1 - L_{\Gamma,s} \right), \quad \text{Re} \, s \gg 1
\]

where \(h\) is a meromorphic function accounting for certain non-exact codings. A weak version of our second main result reads as follows.

**Theorem D** Let \(\Gamma\) be a geometrically finite Fuchsian group which admits a strict transfer operator approach, and let \(\chi\) be a finite-dimensional representation of \(\Gamma\)
with non-expanding cusp monodromy. Then the infinite product (1) for $Z_{\Gamma, \chi}$ extends meromorphically to all of $\mathbb{C}$. The poles of $Z_{\Gamma, \chi}$ are contained in $\frac{1}{2}(d_1 - N_0)$, where $d_1$ denotes the maximal length of a Jordan chain of $\chi(p)$ with eigenvalue $1$, $p \in \Gamma$ parabolic.

In fact, we will prove a stronger version of Theorem D (see Theorem 4.2 below) for which we first show that the infinite product (1) is represented by a twisted version of the transfer operator family:

$$Z_{\Gamma, \chi}(s) = \det (1 - L_{\Gamma, \chi, s}), \quad \Re s \gg 1. \quad (3)$$

Then we will prove that not only the Selberg zeta function $Z_{\Gamma, \chi}$ and the Fredholm determinant $s \mapsto \det (1 - L_{\Gamma, \chi, s})$ admit a meromorphic continuation but also the map

$$s \mapsto L_{\Gamma, \chi, s} \quad (4)$$

itself. In addition, we will provide upper bounds on the order of the poles.

The requirement that $\Gamma$ admit a strict transfer operator approach does not directly ask for the equality (2) but for the existence of a suitable (uniformly expanding) iterated function system (IFS). If the representation $\chi$ is unitary, then establishing (2), the meromorphic continuation of (4) and (the strong version of) Theorem D starting from such an IFS is by now standard, see, e.g., [16,38,50,52].

Hence, the focus in the proof of Theorem D is to accommodate non-unitary representations. For (3), nuclearity of the twisted transfer operators is indispensable. To show this property we will rely on the logarithm law for cusp excursions as well as on the finiteness of the number of cusps of $\Gamma$.

If $\chi$ is the trivial one-dimensional representation, then the standard proof of the meromorphic continuation of (4) takes advantage of the Hurwitz zeta function. For any unitary representations $\chi$, Pohl [50] recently showed that this proof can be adapted by combining a diagonalization of $\chi(p)$, $p \in \Gamma$ parabolic, with the Lerch zeta function. For generic representations $\chi$ with non-expanding cusp monodromy, the endomorphisms $\chi(p)$, $p \in \Gamma$ parabolic, are not necessarily diagonalizable anymore. However, as we will show in Sect. 4 below, a careful use of the Lerch transcendent and its derivatives in combination with Jordan normal forms of $\chi(p)$, $p \in \Gamma$ parabolic, allows us to establish meromorphic continuability.

As an additional result to those already mentioned, we will briefly discuss, in Sect. 7 below, that the Venkov–Zograf factorization formulas also hold for Selberg zeta functions with twists by representations with non-expanding cusp monodromy.

The next natural goal in this line of research is to seek for a spectral and topological characterization of the zeros of the Selberg zeta function $Z_{\Gamma, \chi}$. Answering this question we leave for future research. The first step in this direction is already done in [15], where the convergence of $\chi$-twisted Eisenstein series is addressed.
2 Preliminaries

2.1 Hyperbolic geometry

We use the upper half-plane model of the hyperbolic plane

$$\mathbb{H} := \{ x + iy \mid y > 0 \}$$

endowed with the hyperbolic metric

$$ds^2 = \frac{dx^2 + dy^2}{y^2}. \quad (5)$$

The group of orientation-preserving Riemannian isometries of $\mathbb{H}$ is isomorphic to the group $\text{PSL}_2(\mathbb{R}) = \text{SL}_2(\mathbb{R})/\{ \pm \text{id} \}$, considered as acting by fractional linear transformations on $\mathbb{H}$.

By continuity, this action extends to the geodesic boundary $\partial \mathbb{H}$ of $\mathbb{H}$. If we identify $\partial \mathbb{H}$ with $\hat{\mathbb{R}} := \mathbb{R} \cup \{ \infty \}$ (or $P^1(\mathbb{R})$) in the standard way, then the action on the geodesic closure

$$\overline{\mathbb{H}} = \mathbb{H} \cup \partial \mathbb{H}$$

of $\mathbb{H}$ reads as follows: For $g = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \in \text{PSL}_2(\mathbb{R})$ and $z \in \overline{\mathbb{H}}$ we have

$$g.z := \begin{cases} \infty & \text{if } z = \infty, c = 0, \text{ or } z \neq \infty, cz + d = 0 \\ \frac{a}{c} & \text{if } z = \infty, c \neq 0 \\ \frac{az + b}{cz + d} & \text{otherwise.} \end{cases}$$

For any hyperbolic element $h \in \text{PSL}_2(\mathbb{R})$ (i.e., $|\text{tr}(h)| > 2$) we let $N(h)$ denote the norm of $h$, thus

$$N(h) := \max \left\{ \lambda^2 \mid \lambda \text{ eigenvalue of } h \right\}. \quad (6)$$

2.2 Hyperbolic surfaces/orbifolds and periodic geodesics

Let $\Gamma$ be a Fuchsian group. We denote by $[\Gamma]_h$ the set of $\Gamma$-conjugacy classes of hyperbolic elements in $\Gamma$, and by $[\Gamma]_p$ the subset of $\Gamma$-conjugacy classes of primitive hyperbolic elements. For $h \in \Gamma$ hyperbolic, $[h]$ denotes the corresponding element in $[\Gamma]_h$.

Let

$$X := \Gamma \setminus \mathbb{H}$$

denote the (two-dimensional, connected, real hyperbolic, good) orbifold with fundamental group $\Gamma$. As is well-known, the periodic geodesics on $X$ are in bijection...
A natural isomorphism between periodic geodesics and $[\Gamma]_p$, on which we will rely in Sect. 3 below, is given as follows: To a hyperbolic element $h \in \Gamma$ we assign the geodesic $\gamma_h$ on $H$ for which $\gamma_h(+\infty)$ is the attracting fixed point of $h$, and $\gamma_h(-\infty)$ is the repelling fixed point. Note that $\gamma_h^n = \gamma_h$ for any $n \in \mathbb{N}$. Then we identify an equivalence class $[g] \in [\Gamma]_p$ of a primitive hyperbolic element $g \in \Gamma$ with the periodic geodesic $\hat{\gamma}_g$ on $X$ for which $\gamma_g$ is a representing geodesic on $H$. In this case, the (primitive) length $\ell(\hat{\gamma}_g)$ of $\hat{\gamma}_g$ is given by

$$\ell(\hat{\gamma}_g) = \log N(g).$$

We remark that $\hat{\gamma}_g$ and $N(g)$ do not depend on the chosen representative $g \in \Gamma$ of $[g] \in [\Gamma]_p$. If we consider (as we shall) periodic geodesics on $X$ with length multiplicities, then this identification extends to all of $[\Gamma]_h$.

For any hyperbolic $h \in \Gamma$ we set

$$\ell(h) = \log N(h).$$

### 2.3 Selberg zeta function

Let $\Gamma$ be a geometrically finite Fuchsian group, and let $\chi : \Gamma \to \text{GL}(V)$ be a finite-dimensional representation of $\Gamma$. The Selberg zeta function $Z_{\Gamma,\chi} = Z(\cdot, \Gamma, \chi)$ for $(\Gamma, \chi)$ is formally defined by

$$Z_{\Gamma,\chi}(s) := Z(s, \Gamma, \chi) := \prod_{[g] \in [\Gamma]_p} \prod_{k=0}^{\infty} \det \left( 1 - \chi(g) N(g)^{-s-k} \right),$$

where $1$ denotes the identity map on $V$. The relation to the definition in (1) in the Introduction is explained by (7).

If $\chi$ is the trivial one-dimensional representation $1$, then the infinite product (8) is known to converge for $\text{Re} \, s \gg 1$ (more precisely, for $\text{Re} \, s > \delta$, where $\delta$ denotes the Hausdorff dimension of the limit set of $\Gamma$) [21,45,53,56,57]. As recalled in the Introduction, convergence is more generally known if $\chi$ is unitary—a result we do not take advantage of in this article.

For the elementary Fuchsian groups, the Selberg zeta function is most interesting for the hyperbolic funnel groups, that is, those Fuchsian groups that are generated by a single hyperbolic element. For the other elementary Fuchsian groups, the first product in (8) is void, thus the Selberg zeta function is constant 1 by convention.

### 2.4 Representations

Let $\Gamma$ be a Fuchsian group. A finite-dimensional (real) representation $\chi : \Gamma \to \text{GL}(V)$ is said to have non-expanding cusp monodromy if for each parabolic element $p \in \Gamma$, each eigenvalue of the endomorphism $\chi(p)$ has absolute value 1.
2.5 Further notations

We use $\mathbb{N} = \{1, 2, 3, \ldots\}$ to denote the set of natural numbers (without zero), and let $\mathbb{N}_0 := \mathbb{N} \cup \{0\}$. Further, for functions $f : M \to \mathbb{C}$, $g : M \to \mathbb{R}$ on some set $M$ we write $|f| \ll g$ or $f \ll g$ or $|f(m)| \ll g(m)$ or $f(m) \ll g(m)$ for

$$\exists C > 0 \forall m \in M : |f(m)| \leq C g(m).$$

If the constant $C$ depends on an exterior variable, say $p$, and we need to keep track of this dependence, then we indicate this variable in an subscript, e.g., $\ll_p$. Analogously, we use $\gg$ and $\gg_p$.

3 Convergence of twisted Selberg zeta functions

Throughout this section let $\Gamma$ be a geometrically finite Fuchsian group and $\chi : \Gamma \to \text{GL}(V)$ a finite-dimensional representation with non-expanding cusp monodromy. We will show that the infinite product in (8) converges if Re $s$ is sufficiently large (depending on $\Gamma$ and $\chi$).

**Theorem 3.1** There exists $C = C(\Gamma, \chi) \in \mathbb{R}$ such that the infinite product in (8) converges compactly on the right half-plane $R := \{s \in \mathbb{C} | \text{Re } s > C\}$. Thus, on $R$, it defines a holomorphic function without zeros.

As already mentioned in the Introduction, the proof of Theorem 3.1 is based on bounding the Selberg zeta function $Z_{\Gamma, \chi}$ for the representation $\chi$ by the Selberg zeta function $Z_{\Gamma, 1}$ for the trivial one-dimensional representation $1$, and then taking advantage of the fact that the convergence of $Z_{\Gamma, 1}$ is well-known. More precisely, we find $c \in \mathbb{R}$ such that for all $s \in \mathbb{C}$ with Re $s$ sufficiently large we have

$$\log Z(s, \Gamma, \chi) \ll |\log Z(\text{Re } s - c, \Gamma, 1)|.$$

Proposition 3.2 below is the key step for the proof that the shift in the argument of $Z_{\Gamma, 1}$ is indeed uniform. This proposition makes crucial use of $\chi$ having non-expanding cusp monodromy.

Throughout we fix a norm $\|\cdot\|$ on $\text{GL}(V)$. Since the trace norm is sub-multiplicative, and all norms on $\text{GL}(V)$ are equivalent due to finite-dimensionality, $\|\cdot\|$ is essentially sub-multiplicative. Thus, there exists $C > 0$ such that for all $g, h \in \Gamma$ we have the bound

$$\|\chi(gh)\| \leq C \|\chi(g)\| \cdot \|\chi(h)\|. \tag{9}$$

**Proposition 3.2** There exists $c \geq 0$ such that uniformly for all hyperbolic elements $h \in \Gamma$, we have

$$|\text{tr } \chi(h)| \ll e^{c \ell(h)}, \tag{10}$$
or, equivalently,

\[ |\text{tr} \chi(h)| \ll N(h)^c. \]

Moreover,

\[ \inf_{g \in \Gamma} \|\chi(ghg^{-1})\| \ll N(h)^c. \]

In particular, all implied constants are independent of \( h \).

For elementary Fuchsian groups, Proposition 3.2 is obvious. For the proof of Proposition 3.2 for non-elementary Fuchsian groups we need some preparations. In Sect. 3.1 below we will recall an estimate of word lengths (or rather ‘block’ lengths) of hyperbolic elements in terms of the displacement of the point \( i \) of the hyperbolic plane \( \mathbb{H} \) as provided by Eichler [12] and by Knopp and Sheingorn [30]. This result allows us to bound the number of cusp excursions of a periodic geodesic.

In Sect. 3.2 below we will provide, for every parabolic element \( p \in \Gamma \), a polynomial bound on \( \chi(p^m) \) as \( m \to \infty \) (the logarithm-type law mentioned in the Introduction). In Sect. 3.3 below we will recall the notion of horoballs and their relation to Dirichlet fundamental domains, which obviously is closely related to Siegel sets. For every periodic geodesic \( \hat{\gamma} \) on \( \Gamma \backslash \mathbb{H} \), the latter relation allows us to estimate how deep \( \hat{\gamma} \) travels into the cusp represented by the parabolic element \( p \). This permits us to deduce bounds on the power \( m \) with which \( p^m \) appears in the word/block representation of a hyperbolic element \( h \) associated to \( \hat{\gamma} \). These bounds grow at most polynomial (with uniform degree) in the norm \( N(h) \) of \( h \) or, equivalently, at most exponential in the length \( \ell(\hat{\gamma}) \) of \( \hat{\gamma} \).

### 3.1 Bounds on block lengths

We recall an upper estimate of ‘block’ lengths of hyperbolic elements \( h \) in terms of the displacement of the point \( i \) by \( h \). This estimate was proved by Eichler [12] for Fuchsian lattices and was then generalized by Knopp and Sheingorn [30] to all geometrically finite Fuchsian groups.

Let \( D \) be a Dirichlet fundamental domain for \( \Gamma \), and let \( \mathcal{A} \subseteq \Gamma \) be the set of side-pairing elements of \( D \). We present each element \( g \in \Gamma \) as a word over \( \mathcal{A} \) via the standard Morse coding algorithm induced by \( D \). This means that the presentation \( g = A_1A_2 \cdots A_n \) as a word over \( \mathcal{A} \) is deduced as follows: We fix a point \( x \) in \( D \) and a point \( y \) in \( g.D \) such that the geodesic \( \gamma \) (on \( \mathbb{H} \)) connecting \( x \) and \( y \) stays away from the vertices of \( k.D \) for all \( k \in \Gamma \). Thus, when moving along the geodesic \( \gamma \) from \( x \) to \( y \), one passes through a (unique) sequence of the form

\[ D, A_1.D, A_1A_2.D, \ldots, A_1A_2 \cdots A_n.D = g.D \]

with \( A_1, \ldots, A_n \in \mathcal{A} \). This sequence is indeed independent of the choice of \( x \) and \( y \) as long as the connecting geodesic \( \gamma \) satisfies the condition mentioned. In each step, the
element $A_i \in \mathcal{A}$ is determined by the side of $A_1 \cdots A_{i-1}D$ through which $\gamma$ passes, or equivalently, the side $S_i$ of $D$ through which $A_{i-1}^{-1} \cdots A_1^{-1}\gamma$ passes. It is the unique element $A_i$ in $\Gamma \backslash \{\text{id}\}$ (or, equivalently, in $\mathcal{A}$) such that $A_iS_i$ is again a side of $D$.

We arrange the presentation $g = A_1 \cdots A_n$ as follows into ‘blocks’: Each non-parabolic element $A \in \mathcal{A}$ forms a block on its own, and each locally maximal subword of $A_1 \cdots A_n$ of the form $A^k = A \cdots A$ ($k$-times) with $A \in \mathcal{A}$ being parabolic and $k \in \mathbb{N}$ forms a block. Here, ‘locally maximal’ means that

$$A_1 \cdots A_n = A_1 \cdots A_i A^{k} A_{i+k+1} \cdots A_n$$

with $A_i \neq A \neq A_{i+k+1}$ (thus, $k$ is maximal). Let $b_D(g)$ denote the number of blocks of $g$. Note that $b_D$ may depend on the choice of $D$.

Let $d_H$ denote the metric on $H$ induced by the Riemannian metric (5). By Eichler [12] and Knopp and Sheingorn [30] there exist constants $c_1, c_2 > 0$, possibly depending on the choice of $D$, such that for all hyperbolic elements $h \in \Gamma$, we have

$$b_D(h) \leq c_1 \cdot d_H(i, h) + c_2.$$ (11)

### 3.2 Growth of norms

We show that for every parabolic element $p \in \Gamma$, the norm of $\chi(p^m)$ grows at most polynomially as $m \to \infty$. The degree of the polynomial is uniformly (over all parabolic elements) bounded by $\dim V - 1$.

**Lemma 3.3** Let $p \in \Gamma$ be parabolic, $d_p \in \mathbb{N}$, and suppose that the lengths of the Jordan chains of $\chi(p)$ are bounded by $d_p$. Then

$$\|\chi(p^m)\| \ll_p m^{d_p-1}$$

for all $m \in \mathbb{N}$.

**Proof** We fix a basis of $V$ with respect to which $\chi(p)$ is represented in Jordan normal form. Let $r \in \mathbb{N}$ be the number of Jordan blocks, $k_1, \ldots, k_r \in \mathbb{N}$ the lengths of the Jordan chains, and $\lambda_1, \ldots, \lambda_r \in \mathbb{C}$ the corresponding eigenvalues of $\chi(p)$. Then

$$\|\chi(p^m)\| \ll_p \max_{j=1}^{r} \max_{\ell=0}^{k_j-1} \left|\lambda_j\right|^{m-\ell} \left(\frac{m}{\ell}\right).$$

By hypothesis, $k_j \leq d_p$ for $j = 1, \ldots, r$. Further, since $\chi$ has non-expanding cusp monodromy, $|\lambda_j| = 1$ for $j = 1, \ldots, r$. Thus

$$\max_{j=1}^{r} \max_{\ell=0}^{k_j-1} \left|\lambda_j\right|^{m-\ell} \left(\frac{m}{\ell}\right) \leq \max_{l=0}^{d_p-1} \left(\frac{m}{l}\right) = \begin{cases} \left(\frac{m}{d_p-1}\right) & \text{for } d_p - 1 < \left\lfloor \frac{m}{2} \right\rfloor, \\ \left(\frac{m}{d_p-1}\right) & \text{for } d_p - 1 \geq \left\lfloor \frac{m}{2} \right\rfloor. \end{cases}$$
Hence,
\[
\max_{j=1}^r \max_{\ell=0}^{k_j-1} \left\{ |\lambda_j|^{m-\ell} \left( \frac{m}{\ell} \right) \right\} \ll m^{d_{p}-1}.
\]

This completes the proof. \(\square\)

### 3.3 Return bounds

Recall that \(d_\mathbb{H}\) denotes the metric on \(\mathbb{H}\) induced by the Riemannian metric (5), and let \(\gamma\) be a geodesic on \(\mathbb{H}\). Then the Buseman function \(b_\gamma: \mathbb{H} \to \mathbb{R}\) is defined by

\[
b_\gamma(z) := \lim_{t \to \infty} \left( t - d_\mathbb{H}(z, \gamma(t)) \right).
\]

Let \(c \in \partial \mathbb{H}\) and \(r \in \mathbb{R}\). The horoball \(B(c, r)\) centered at \(c\) with ‘radius’ \(r\) is the sublevel set

\[
B(c, r) := b_\gamma^{-1}([r, \infty)),
\]

where \(\eta\) is any geodesic on \(\mathbb{H}\) with \(\eta(\infty) = c\). If \(c \in \mathbb{H}\) is cuspidal with respect to \(\Gamma\), i.e., \(c\) is fixed by a primitive parabolic element \(p\) in \(\Gamma\), then there exists a radius \(r = r(c)\) such that whenever \(z, w \in B(c, r)\) are \(\Gamma\)-equivalent, hence \(g.z = w\) for some \(g \in \Gamma\), then \(g = p^m\) for some \(m \in \mathbb{Z}\). Even more, if \(D\) is a Dirichlet fundamental domain for \(\Gamma\), \(\mathcal{A}\) is the set of side-pairing elements, and \(\mathcal{P} \subseteq \mathcal{A}\) its subset of the parabolic elements among the side-pairing elements, then there exists \(r_0 \in \mathbb{R}\) such that for each \(p \in \mathcal{P}\),

\[
B(c_p, r_0) \subseteq \bigcup_{m \in \mathbb{Z}} p^m.D,
\]

where \(c_p\) denotes the fixed point of \(p\) (see, e.g., [2, Section 9.4], in particular [2, Theorem 9.4.5]). We call each such \(r_0 \in \mathbb{R}\) a return bound for \(D\).

### 3.4 Proof of Proposition 3.2

Let \(D\) be a Dirichlet fundamental domain for \(\Gamma\) which we fix once and for all. In the following, all of the implied constants as well as the numerically unspecified constants may depend on the choice of \(D\) but they are independent of any specially considered element in \(\Gamma\).

Let \(\mathcal{A}\) denote the set of side-pairing elements of \(D\), let

\[
K := \max_{A \in \mathcal{A}} \| \chi(A) \|
\]
and recall the definition of block presentations and the block counting function

\[ b := b_D \]

from Sect. 3.1. Further let \( D_0 \) denote the (compact) subset of \( \overline{D} \) which represents the compact core of \( \Gamma \backslash \mathbb{H} \) (or, more generally, let \( D_0 \) be any compact subset of \( \overline{D} \) which represents a compact subset of \( \Gamma \backslash \mathbb{H} \) that is intersected by all periodic geodesics).

Since every periodic geodesic on \( \Gamma \backslash \mathbb{H} \) intersects the compact core, every \( \Gamma \)-conjugacy class of hyperbolic elements in \( \Gamma \) contains at least one representative \( h \) whose associated geodesic \( \gamma_h \) on \( \mathbb{H} \) intersects \( D_0 \). It suffices to prove (10) for one such representative \( h \) out of each \( \Gamma \)-conjugacy class of hyperbolic elements since \( \text{tr} \circ \chi \) as well as the length function \( \ell \) are constant on these conjugacy classes.

Let \( h \in \Gamma \) be a hyperbolic element such that its associated geodesic \( \gamma_h \) on \( \mathbb{H} \) intersects \( D_0 \). Suppose that

\[ h = h_1 \cdots h_{b(h)} \]

is its block presentation. Let

\[ J_p := \{ j \in \{1, \ldots, b(h)\} \mid h_j \text{ parabolic} \} \]

be the set of indices corresponding to the parabolic blocks, and set

\[ J_n := \{1, \ldots, b(h)\} \setminus J_p. \]

Further set

\[ b_p := \#J_p \quad \text{and} \quad b_n := \#J_n, \]

so that \( b_p + b_n = b(h) \). Then

\[
|\text{tr} \chi(h)| \ll \|\chi(h)\| \leq C^{b(h)-1} \prod_{j=1}^{b(h)} \|\chi(h_j)\|
\]

\[
= C^{b(h)-1} \left( \prod_{i \in J_n} \|\chi(h_i)\| \right) \left( \prod_{j \in J_p} \|\chi(h_j)\| \right)
\]

(12)

with \( C \) as in (9). For each \( i \in J_n \), the block \( h_i \) consists of a single element in \( \mathcal{A} \). Hence

\[
\prod_{i \in J_n} \|\chi(h_i)\| \leq K^{b_n}. \quad (13)
\]

For each \( k \in J_p \), the block \( h_j \) equals \( p_j^{m_j} \) for some parabolic element \( p_j \in \mathcal{A} \) and \( m_j \in \mathbb{N} \). Let \( d_0 \) denote the maximal length of a Jordan chain of \( \chi(p), p \in \mathcal{A} \) parabolic.
Further let $\tilde{C} > 0$ be an upper bound for the implied constants in Lemma 3.3, when applied to the finitely many parabolic elements in $A$. Then Lemma 3.3 yields

$$\prod_{j \in J_p} \| \chi(h_j) \| \leq \tilde{C} b_p \prod_{j \in J_p} m_j^{d_0 - 1}. \quad (14)$$

In order to bound $m_j$ for $j \in J_p$ we fix a point $z \in \gamma_h(\mathbb{R}) \cap D_0$. Without loss of generality, we may assume that $z$ serves to deduce the presentation of $h$ as a word over $A$ as in Sect. 3.1. Thus, the geodesic segment $c_h$ connecting $z$ and $h.z$ is a segment of $\gamma_h$, and

$$d_{\mathbb{H}}(z, h.z) = \ell(h) = \ell(c_h). \quad (15)$$

Here, $\ell(c_h)$ denotes the length of $c_h$.

Let $j \in J_p$. Let $b_j$ denote the fixed point of the parabolic element $p_j$ and let $\gamma_j := h_{j-1}^{-1} \cdots h_1^{-1} \gamma_h$. Since $h_j = p_j^{m_j}$, the geodesic $\gamma_j$ passed through $A.D$ for some $A \in A \setminus \{p_j\}$ before entering $D$, then passes through $D, p_j.D, \ldots, p_j^{m_j}.D$ and finally enters into $p_j^{m_j}A'.D$ for some $A' \in A \setminus \{p_j\}$. Set

$$D_j := \bigcup_{k=0}^{m_j} p_j^k.D$$

and let

$$t_j := \mu_{\text{Leb}} \left( \{ t \in \mathbb{R} \mid \gamma_j(t) \in D_j \} \right)$$

be the time (or, equivalently, the length) that $\gamma_j$ spends in $D_j$. Here, $\mu_{\text{Leb}}$ denotes the Lebesgue measure on $\mathbb{R}$ (normalized such that $\mu_{\text{Leb}}([0, 1]) = 1$). Let $S_1, S_2$ be the two sides of $D$ which meet at $b_j$, and suppose that $\gamma_j$ exits $D$ through $S_2$. Then, for $k = 1, \ldots, m_j - 1$, it enters $p_j^k.D$ through $p_j^k.S_1 = p_j^{k-1}.S_2$ and exits through $p_j^k.S_2$. Suppose that

$$a_j := \inf_{k=1}^{m_j-1} \mu_{\text{Leb}} \left( \{ t \in \mathbb{R} \mid \gamma_j(t) \in p_j^k.D \} \right)$$

denotes the minimal time that $\gamma_j$ spends in any of the $p_j^k.D, k = 1, \ldots, m_j - 1$. Then

$$m_j - 1 \leq \frac{t_j}{a_j}.$$ 

In order to provide a lower bound on $a_j$, let $r_0 \in \mathbb{R}$ be a return bound for $D$ (cf. Sect. 3.3) and recall that

$$B(b_j, r_0) \subseteq \bigcup_{n \in \mathbb{Z}} p_j^n.D.$$
Obviously, the time that $\gamma_j$ spends in $B(b_j, r_0)$ is bounded by $t_j$. Therefore, $\gamma_j$ does not enter $B(b_j, r_0 + t_j)$. Let

$$H := D \cap B(b_j, r_0) \setminus B(b_j, r_0 + t_j)$$

and let $C$ denote the set of geodesics $\gamma$ on $\mathbb{H}$ that enter $H$ through $S_1$ and exit through $S_2$. Then $a_j$ is bounded below by

$$a'_j := \inf_{\gamma \in C} \mu_{\text{Leb}}((t \in \mathbb{R} \mid \gamma(t) \in H)).$$

Let $w_j$ be the cusp width of $D$ at $b_j$. By elementary hyperbolic geometry,

$$a'_j = \text{arcosh} \left(1 + \frac{w_j}{2e^{2t_j}}e^{-2t_j}\right) \gg e^{-3t_j}.$$

Therefore,

$$m_j \leq \frac{t_j}{a_j} + 1 \ll t_j e^{3t_j} + 1 \ll e^{4t_j}.$$

Since

$$\sum_{j \in J_p} t_j \leq \ell(c_h) = \ell(h),$$

we find

$$\prod_{j \in J_p} m_j^{d_0-1} \ll e^{A(d_0-1)\ell(h)}.$$  \hspace{1cm} (16)

The combination of (12), (13), (14) and (16) yields

$$C|\text{tr } \chi(h)| \ll C\|\chi(h)\| \leq C_1^{b(h)} e^{A(d_0-1)\ell(h)}$$  \hspace{1cm} (17)

for some $C_1 > 0$. By (11),

$$C_1^{b(h)} \ll e^{c_1 \cdot d_\mathbb{H}(i, h.i)}$$

for some $c_1 > 0$. The triangle inequality, the left-invariance of the metric $d_\mathbb{H}$, and the compactness of $D_0$ imply the existence of $c_2 \geq 0$ such that

$$d_\mathbb{H}(i, h.i) \leq d_\mathbb{H}(i, z) + d_\mathbb{H}(z, h.z) + d_\mathbb{H}(h.z, h.i) \leq 2d_\mathbb{H}(i, z) + d_\mathbb{H}(z, h.z) \leq c_2 + \ell(h).$$  \hspace{1cm} (18)

Hence there exists $c \geq 0$ such that

$$|\text{tr } \chi(h)| \ll \|\chi(h)\| \ll e^{c\ell(h)}.$$

This proves Proposition 3.2. \hfill \square
3.5 Proof of Theorem 3.1

For any hyperbolic element \( g \in \Gamma \) let \( m(g) \in \mathbb{N} \) denote the unique number (‘multiplicity’) such that \( g = g_0^{m(g)} \) for a primitive hyperbolic element \( g_0 \in \Gamma \).

By Proposition 3.2 we find \( c \geq 0 \) (which we fix throughout) such that uniformly for all hyperbolic elements \( h \in \Gamma \) we have

\[
|\text{tr} \chi(h)| \ll N(h)^c \tag{19}
\]

and

\[
\inf_{k \in \Gamma} \|\chi(khk^{-1})\| \ll N(h)^c. \tag{20}\]

For \( s \in \mathbb{C}, \text{Re} \ s > c \), the bound (20) yields that for each \([g] \in [\Gamma]_p\) there exists a representative \( g \in \Gamma \) such that

\[
\|\chi(g)N(g)^{-s}\| < 1.
\]

It follows that, for all \( k \in \mathbb{N}_0 \), the equality

\[
\text{tr} \log \left( 1 - \chi(g)N(g)^{-(s+k)} \right) = -\sum_{n=1}^{\infty} \frac{N(g^n)^{-(s+k)}}{n} \text{tr} \chi(g^n)
\]

is valid and the right hand side’s series is convergent. Therefore, for \( \text{Re} \ s > c \), we have the formal identity

\[
\log Z(s, \Gamma, \chi) = -\sum_{[g] \in [\Gamma]_h} \frac{1}{m(g)} \frac{N(g)^{-s}}{1 - N(g)^{-1}} \text{tr} \chi(g) \tag{21}
\]

and the convergence of the infinite product (8) is equivalent to the convergence of either side of (21).

Recall that \( \mathbf{1} \) denotes the trivial one-dimensional representation of \( \Gamma \). Combining (21) and (19) we find

\[
|\log Z(s, \Gamma, \chi)| \leq \sum_{[g] \in [\Gamma]_h} \frac{1}{m(g)} \frac{N(g)^{-\text{Re} s}}{1 - N(g)^{-1}} |\text{tr} \chi(g)|
\]

\[
\ll \sum_{[g] \in [\Gamma]_h} \frac{1}{m(g)} \frac{N(g)^{-(\text{Re} s - c)}}{1 - N(g)^{-1}}
\]

\[
= |\log Z(\text{Re} s - c, \Gamma, \mathbf{1})|.
\]

The map \( s \mapsto \log Z(s - c, \Gamma, \mathbf{1}) \) converges compactly on a right half-plane in \( \mathbb{C} \) (see Sect. 2.3). More precisely, it converges compactly on \( \{\text{Re} s > \delta + c\} \), where \( \delta \) is the Hausdorff dimension of the limit set of \( \Gamma \). This completes the proof. \( \square \)
3.6 A side-result

A variation of the proof of Proposition 3.2 shows an interesting relation between the norm of $\| \chi(h) \|$ and the displacement $d_H(z, h.z)$ for all hyperbolic elements $h \in \Gamma$. Even though we do not need this result in this article, we record it here for purposes of reference. This result is crucial for the study of $\chi$-twisted Eisenstein series [15].

**Corollary 3.4** There exists $c \geq 0$ and a continuous map $\alpha : \mathbb{H} \to \mathbb{R}$ such that uniformly for all hyperbolic elements $h \in \Gamma$ and all $z \in \mathbb{H}$, we have

$$\| \chi(h) \| \leq \alpha(z) e^{c \cdot d_H(z, h.z)}.$$

**Proof** We use the notation from the proof of Proposition 3.2. Without loss of generality we suppose that the fundamental domain is chosen such that $i \in D$ (we could pick any other generic point $z_0$ and arrange $D$ around $z_0$).

Let $h \in \Gamma$ be an arbitrary hyperbolic element. We do not suppose that $\gamma_h$ intersects the compact core $D_0$. We suppose that $i$ is the point to deduce the block decomposition of $h$ as in Sect. 3.1. In the unlikely case that the geodesic connecting $i$ and $h.i$ passes through a vertex of a $\Gamma$-translate of $D$, we use a slight perturbation of $i$. Any compact perturbation does not change the nature of the results.

We write $h$ in its block decomposition and proceed to estimate $\| \chi(h) \|$ as in the proof of Proposition 3.2 until (14). In order to estimate the $m_j$’s we let $c_h$ be the geodesic segment connecting $i$ and $h.i$. Note that we do not request that $i$ is in $\gamma_h(\mathbb{R})$. This is a major difference to the proof of Proposition 3.2. Then (15) becomes

$$d_H(i, h.i) = \ell(c_h).$$

We continue to proceed as in the proof of Proposition 3.2 until (16), which now reads

$$\prod_{j \in J_p} m_j^{d_0-1} \ll e^{4(d_0-1)d_H(i.h.i)}.$$

It follows that

$$\| \chi(h) \| \ll e^{c \cdot d_H(i.h.i)},$$

where

$$c = c_1 + 4(d_0 - 1).$$

Let $z \in \mathbb{H}$. Using (18) shows

$$\| \chi(h) \| \ll e^{2d_H(i,z)} e^{c \cdot d_H(z, h.z)}.$$

This completes the proof. □
4 Meromorphic continuation of twisted Selberg zeta functions

Throughout this section let $\Gamma$ be a geometrically finite Fuchsian group and let $\chi : \Gamma \to \text{GL}(V)$ be a finite-dimensional representation with non-expanding cusp monodromy. In this section we will show that under the assumption that $\Gamma$ admits a so-called strict transfer operator approach (defined in Sect. 4.2 below), the Selberg zeta function $Z_{\Gamma, \chi}$ admits a meromorphic continuation to all of $\mathbb{C}$. More precisely, we will show the stronger statement that $Z_{\Gamma, \chi}$ equals the Fredholm determinant of a family of twisted transfer operators, and that this family admits a meromorphic continuation to all of $\mathbb{C}$. In addition, we will provide upper bounds on the order of poles in the meromorphic continuation. For precise statements we refer to Theorem 4.2 below.

In order to define the notion of a strict transfer operator approach we need a few preparations.

4.1 Geometry

We consider the Riemann sphere $\hat{\mathbb{C}} = \mathbb{C} \cup \{\infty\}$ as a complex manifold with the standard complex structure which is determined by two charts, the first of which is given by the identity on $\mathbb{C}$, and the second is

$$\hat{\mathbb{C}} \setminus \{0\} \to \mathbb{C}, \quad z \mapsto \begin{cases} 0 & \text{if } z = \infty \\ \frac{1}{z} & \text{otherwise.} \end{cases}$$

We call a subset of $\hat{\mathbb{C}}$ bounded if its image in at least one of these charts is bounded as a subset of $\mathbb{C}$ (endowed with the euclidean metric). Note that this notion of boundedness does not coincide with boundedness with respect to the chordal metric on $\hat{\mathbb{C}}$. Note further that a bounded subset of $\hat{\mathbb{C}}$ is in particular contained in the domain of a chart, and hence $\hat{\mathbb{C}}$ itself is not a bounded set.

Likewise we consider $\hat{\mathbb{R}} = \mathbb{R} \cup \{\infty\}$ as a real manifold whose manifold structure is determined by the restriction of the two charts for $\hat{\mathbb{C}}$ to $\mathbb{R}$. We call a subset of $\hat{\mathbb{R}}$ an interval if its image in at least one of these charts is an interval in $\mathbb{R}$. According to this definition, the empty set and singletons are intervals. The whole manifold $\hat{\mathbb{R}}$ itself however is not considered to be an interval because it is not contained in the domain of a single chart map.

Let $\hat{\mathcal{V}}_f$ denote the set of geodesics on $X := \Gamma \backslash \mathbb{H}$ which converge to a cusp or funnel in forward time direction. In other words, $\hat{\mathcal{V}}_f$ is the set of geodesics which, in forward time direction, eventually leave any compact set forever. Let $\mathcal{V}_f$ denote the (maximal) set of geodesics on $\mathbb{H}$ which represent the geodesics in $\hat{\mathcal{V}}_f$, and let

$$\text{bd} := \text{bd}(\Gamma) := \{\gamma(\infty) \mid \gamma \in \mathcal{V}_f\}$$

denote the ‘boundary part’ of $X$ of the geodesic boundary $\hat{\mathbb{R}}$ of $\mathbb{H}$.

We provide a more geometrical description of the set $\text{bd} = \text{bd}(\Gamma)$: Suppose that $\mathcal{F}$ is a fundamental domain for $\Gamma$ in $\mathbb{H}$, e.g., a Dirichlet fundamental domain, and consider the closure $\overline{\mathcal{F}}$ of $\mathcal{F}$ in the geodesic closure
\[ \mathbb{H} = \mathbb{H} \cup \hat{\mathbb{R}} \]

of \( \mathbb{H} \) (which is identical to the closure of \( \mathbb{H} \) in \( \hat{\mathbb{C}} \)). Then

\[ \text{bd} = \Gamma.(\mathcal{F} \cap \hat{\mathbb{R}}). \]

For any subset \( I \) of \( \hat{\mathbb{R}} \) we set

\[ I_{\text{st}} := I \setminus \text{bd}. \]

The subscript ‘st’ abbreviates ‘strong’. It alludes to the notion of a ‘strong cross section’ for the geodesic flow on \( \Gamma \setminus \mathbb{H} \). Such cross sections are a good source for tuples that satisfy the notion of a ‘strict transfer operator approach’ which we will define in the next section (cf. [47]).

4.2 Strict transfer operator approaches

We say that \( \Gamma \) admits a strict transfer operator approach if there exists a tuple

\[ S := (A, (I_a)_{a \in A}, (P_{a,b})_{a,b \in A}, (C_{a,b})_{a,b \in A}, ((g_p)_{p \in P_{a,b}})_{a,b \in A}) \]  \hspace{1em} (22)

consisting of

- a finite (index) set \( A \),
- a family \( (I_a)_{a \in A} \) of (not necessarily disjoint) intervals in \( \hat{\mathbb{R}} \),
- for all \( a, b \in A \) a finite (possibly empty) set \( P_{a,b} \) of parabolic elements in \( \Gamma \) and a finite (possibly empty) subset \( C_{a,b} \) of \( \Gamma \), and
- for all \( a, b \in A \) and all \( p \in P_{a,b} \) an element \( g_p \in \Gamma \) (which may be the identity) which satisfies Properties 1–5 below.

**Property 1**

(i) For all \( a, b \in A \), all \( p \in P_{a,b} \) and \( n \in \mathbb{N} \) we have \( p^{-n}g_p^{-1}I_{a,\text{st}} \subseteq I_{b,\text{st}} \).

Further, \( p^n \not\in P_{a,b} \) for all \( n \geq 2 \).

(ii) For all \( a, b \in A \) and all \( g \in C_{a,b} \) we have \( g^{-1}I_{a,\text{st}} \subseteq I_{b,\text{st}} \).

(iii) For all \( b \in A \) the sets in the family

\[ \left\{ g^{-1}I_{a,\text{st}} \mid a \in A, \ g \in C_{a,b} \right\} \cup \left\{ p^{-n}g_p^{-1}I_{a,\text{st}} \mid a \in A, \ p \in P_{a,b}, \ n \in \mathbb{N} \right\} \]

are pairwise disjoint, and

\[ I_{b,\text{st}} = \bigcup \left\{ g^{-1}I_{a,\text{st}} \mid a \in A, \ g \in C_{a,b} \right\} \]

\[ \cup \bigcup \left\{ p^{-n}g_p^{-1}I_{a,\text{st}} \mid a \in A, \ p \in P_{a,b}, \ n \in \mathbb{N} \right\}. \]

Property 1 induces a discrete dynamical system \((D_{\text{st}}, F)\), where

\[ D_{\text{st}} := \bigcup_{a \in A} I_{a,\text{st}} \times \{a\}, \]  \hspace{1em} (23)
and $F$ is determined by the submaps

$$g^{-1}.I_{a,\text{st}} \times \{b\} \to I_{a,\text{st}} \times \{a\}, \quad (x, b) \mapsto (g.x, a)$$ (24)

and

$$p^{-n}g^{-1}.I_{a,\text{st}} \times \{b\} \to I_{a,\text{st}} \times \{a\}, \quad (x, b) \mapsto (g_p p^n.x, a)$$ (25)

for $a, b \in A, g \in C_{a,b}, p \in P_{a,b}$ and $n \in \mathbb{N}$.

For $n \in \mathbb{N}$ let $\text{Per}_n$ (‘periodic’) denote the subset of $\Gamma$ which consists of all $g \in \Gamma$ for which there exists $a \in A$ such that

$$g^{-1}.I_{a,\text{st}} \times \{a\} \to I_{a,\text{st}} \times \{a\}, \quad (x, a) \mapsto (g.x, a)$$ (26)

is a submap of $F^n$.

**Property 2** The union

$$\text{Per} := \bigcup_{n \in \mathbb{N}} \text{Per}_n$$

is disjoint.

For $n \in \mathbb{N}$ and $h \in \text{Per}_n$ we say that $w(h) := n$ is the word length of $h$. Property 2 yields that the word length is well-defined.

Recall that $[\Gamma]_h$ denotes the set of $\Gamma$-conjugacy classes of hyperbolic elements in $\Gamma$.

**Property 3** (i) All elements of $\text{Per}$ are hyperbolic.

(ii) If $h \in \text{Per}$, then its primitive $h_0$ is also contained in $\text{Per}$.

(iii) For each $[g] \in [\Gamma]_h$ there exists a unique $n \in \mathbb{N}$ such that there exists $h \in \text{Per}_n$ such that $h$ represents $[g]$.

Suppose that $[g] \in [\Gamma]_h$ is represented by $h$ for some $h \in \text{Per}_n$ for some $n \in \mathbb{N}$. Let $m := m(h) \in \mathbb{N}$ denote the unique number such that $h = h_0^m$ for a primitive hyperbolic element $h_0 \in \Gamma$, and let $p(h) := n/m$. Then we set $w(g) := w(h), m(g) := m(h)$ and $p(g) := p(h)$. By Property 3 these values are well-defined.

**Property 4** For $[g] \in [\Gamma]_h$, there are exactly $p(g)$ distinct elements $h \in \text{Per}_{w(g)}$ such that $h$ represents $[g]$.

**Property 5** There exists a family $(\mathcal{E}_a)_{a \in A}$ of open, bounded, connected, simply connected sets in $\hat{\mathbb{C}}$ such that

(i) for all $a \in A$ we have

$$\overline{T_{a,\text{st}}} \subseteq \mathcal{E}_a,$$
(ii) there exists $\xi \in \text{PSL}_2(\mathbb{R})$ such that for all $a \in A$ we have $\xi, \overline{E}_a \subseteq \mathbb{C}$, and for all $b \in A$ and all $g \in C_{a,b}$ we have

$$g^{\xi^{-1}} \cdot \infty \notin \overline{E}_a,$$

(iii) for all $a, b \in A$ and all $g \in C_{a,b}$ we have

$$g^{-1} \cdot \overline{E}_a \subseteq \overline{E}_b,$$

(iv) for all $a, b \in A$, all $p \in P_{a,b}$ there exists a compact subset $K_{a,b,p}$ of $\hat{\mathbb{C}}$ such that for all $n \in \mathbb{N}$ we have

$$p^{-n} g_p^{-1} \cdot \overline{E}_a \subseteq K_{a,b,p} \subseteq \overline{E}_b,$$

(v) for all $a, b \in A$ and all $p \in P_{a,b}$, the set $g_p^{-1} \cdot \overline{E}_a$ does not contain the fixed point of $p$.

If the tuple $\mathcal{S}$ in (22) satisfies Properties 1–5, then we call it a structure tuple for a strict transfer operator approach. The existence of the element $\xi$ in Property 5(ii) allows us to find a chart in the manifold atlas of $\hat{\mathbb{C}}$ which contains all the sets $\overline{E}_a$, $a \in A$. Without loss of generality we assume throughout that $\xi = \text{id}$. Moreover, to avoid some technical discussions, we assume without loss of generality that for all $a \in A$,

$$\text{Re} \overline{E}_a \subseteq \overline{E}_a.$$

As mentioned in the Introduction, the number of geometrically finite Fuchsian groups for which strict transfer operator approaches are established is increasing [1,3, 4,6,16,17,37,41,46–48,50,54]. The hyperbolic funnel groups are (as Fuchsian Schottky groups) among these Fuchsian groups. It is most likely that at least all Fuchsian groups with cusps admit such strict transfer operator approaches, and that our approach can be extended to non-strict transfer operator approaches which then will allow us to cover an even larger class of Fuchsian groups, presumably all geometrically finite ones. In Example 4.1 below we will provide an explicit example for a strict transfer operator approach for the modular group $\text{PSL}_2(\mathbb{Z})$.

For all Fuchsian groups already considered, a strict transfer operator approach can be found for which each of the elements $g_p$ in (22) is the identity. We here allow arbitrary $g_p \in \Gamma$ in order to keep the setting more flexible.

### 4.3 Transfer operators and Banach spaces

From now on we suppose that $\Gamma$ admits a strict transfer operator approach. Let

$$\mathcal{S} := (A, (I_a)_{a \in A}, (P_{a,b})_{a,b \in A}, (C_{a,b})_{a,b \in A}, ((g_p)_{p \in P_{a,b}})_{a,b \in A})$$
be a structure tuple, and let \((D_{st}, F)\) denote the induced discrete dynamical system [(see 23–25)].

As mentioned above, in order to prove the meromorphic continuability of twisted Selberg zeta functions, we will use transfer operators to represent these zeta functions. In this section we will provide the transfer operators associated to the structure tuple and the discrete dynamical system.

For the definition of these transfer operators we will proceed in a two-step process. In the first step we will define these operators only formally, considering them as (formal) operators on the space of functions \(D_{st} \rightarrow V\). This way, we have a clear motivation for the defining expressions of the transfer operators. One immediately sees that these operators are not well-defined on the space of all functions \(D_{st} \rightarrow V\) due to convergence problems. One also sees that there is certain freedom for the choice of the function spaces on which the transfer operators act as actual operators. Therefore, in the second step we will define the Banach spaces which we will use as domains for the transfer operators and will show that (and in which way) the transfer operators constitute well-defined operators on these Banach spaces. See Theorem 4.2 below.

Let

\[
\text{Fct}(D_{st}; V) := \{ f : D_{st} \rightarrow V \text{ function} \}
\]

denote the space of all \(V\)-valued functions on \(D_{st}\). Throughout we fix a norm \(\| \cdot \|\) on \(V\) and an essentially sub-multiplicative norm \(\| \cdot \|\) on \(\text{GL}(V)\). For \(s \in \mathbb{C}\), subsets \(U \subseteq \hat{\mathbb{C}}\), functions \(\psi : U \rightarrow V\), any \(g \in \Gamma\) and any \(z \in \mathbb{C}\) we set

\[
\alpha_s(g^{-1})\psi(z) := \left( g'(z) \right)^s \chi(g^{-1})\psi(g.z),
\]

whenever this definition makes sense. Throughout we use the standard logarithm (i.e., holomorphic on the slit plane \(\mathbb{C} \smallsetminus (-\infty, 0]\)) for the definition of complex powers.

The transfer operator with parameter \(s \in \mathbb{C}\) associated to the structure tuple \(\mathcal{S}\) is formally the operator

\[
\mathcal{L}_s : \text{Fct}(D_{st}; V) \rightarrow \text{Fct}(D_{st}; V)
\]

given by

\[
\left( \mathcal{L}_s f \right)(z) := \sum_{w \in F^{-1}(z)} |F'(w)|^{-s} f(w)
\]

for all \(f \in \text{Fct}(D_{st}; V)\), \(z \in D_{st}\). We note that the space \(D_{st}\) is dense in itself, and hence the derivatives are well-defined.

We identify \(\text{Fct}(D_{st}; V)\) with the space

\[
\{ (f_a)_{a \in A} \mid f_a : I_{a, st} \rightarrow V \text{ for all } a \in A \}
\]

by mapping \(f \in \text{Fct}(D_{st}; V)\) to the function vector \((f_a)_{a \in A}\) determined by

\[
\forall a \in A \forall x \in I_{a, st} : f_a(x) = f(x, a).
\]


With respect to this identification, the transfer operator $\mathcal{L}_s$ is (formally) given by

$$\mathcal{L}_s := (L_{s,a,b})_{a,b \in A}$$

where

$$L_{s,a,b} := \sum_{g \in \mathcal{C}_{a,b}} \alpha_s(g) + \sum_{p \in \mathcal{P}_{a,b}} \sum_{n \in \mathbb{N}} \alpha_s(g_p p^n). \quad (28)$$

Thus, if $f \in \text{Fct}(D_{st}; V)$ and $\tilde{f} = \mathcal{L}_s f$ and $f$ is identified with $(f_a)_{a \in A}$ and $\tilde{f}$ is identified with $(\tilde{f}_a)_{a \in A}$, then

$$\tilde{f}_a = \sum_{b \in A} L_{s,a,b} f_b$$

for all $a \in A$. To find the specific expression (28) of the 'matrix coefficients of the transfer operator' we used the specific form of the submaps of $F$ from (24) and (25).

With domain and range as stated in (27), the transfer operator $\mathcal{L}_s$ is not well-defined as the involved infinite sums do not converge for all elements in $\text{Fct}(D_{st}; V)$. We now provide a domain on which $\mathcal{L}_s$ constitutes a well-defined operator (at least for a certain range of $s$; see Theorem 4.2 below for the proof). To that end we set

$$B(D; V) := \{ \varphi : D \to V \text{ continuous, } \varphi|_| D \text{ holomorphic} \}$$

for any open subset $D \subseteq \mathbb{C}$. We endow $B(D; V)$ with the supremum norm and its norm topology, which turns $B(D; V)$ into a Banach space.

For any family $\mathcal{E}_A = (\mathcal{E}_a)_{a \in A}$ as provided by Property 5 we set

$$B(\mathcal{E}_A; V) := \bigoplus_{a \in A} B(\mathcal{E}_a; V)$$

and endow it with the structure of a Banach space via the product space structure.

The results in this paper do not depend on the choice of the family $\mathcal{E}_A$. This family serves to provide a thickening of the intervals $I_a, a \in A$, into the complex plane such that the relations, in particular the mapping properties, between the intervals are preserved and the operators in (28) are indeed well-defined when acting on $B(\mathcal{E}_A; V)$. To indicate that the essential part of the domain of definitions for the function vectors $(f_a)_{a \in A}$ is the family of intervals $I = (I_a)_{a \in A}$ rather than their thickening, we use the notation

$$B(I; V) := B(\mathcal{E}_A; V) \quad (29)$$

for this Banach space with the understanding that we may choose any family $\mathcal{E}_A$ satisfying Property 5 (the same family for all appearances of $B(I; V)$).

For the convenience of the reader we provide an explicit example of a strict transfer operator approach.

**Example 4.1** Mayer [38,39] conducted the seminal transfer operator approach to a Selberg zeta function, namely for the modular surface $\text{PSL}_2(\mathbb{Z}) \backslash \mathbb{H}$. The discrete
dynamical system his transfer operator is associated to is the Gauss map, which is intimately related to continued fraction expansions and a discretization of the billiard flow with Neumann boundary value conditions on the triangle surface $\text{PGL}_2(\mathbb{Z}) \setminus \mathbb{H}$.

This discretization essentially provides a strict transfer operator approach. However in this setting the discrete group is $\text{PGL}_2(\mathbb{Z})$, which is not Fuchsian. Thus, strictly speaking, Mayer’s transfer operator is not covered by our construction. In turn, this example shows that our construction easily generalizes to discrete subgroups of $\text{PGL}_2(\mathbb{R})$ that are not Fuchsian. To avoid additional technicalities we decided to omit this generalization in this article.

Nevertheless, the Gauss map is also closely related to a discretization of the geodesic flow on the modular surface $\text{PSL}_2(\mathbb{Z}) \setminus \mathbb{H}$. The arising discrete dynamical system and the associated transfer operators have already been discussed in detail in [41,47], for which reason we are here rather brief and only show that the underlying system is a strict transfer operator approach. We use the notation from the previous sections.

Let $\Gamma = \text{PSL}_2(\mathbb{Z})$. The boundary part of $\text{PSL}_2(\mathbb{Z}) \setminus \mathbb{H}$ is

$$\text{bd} = \text{bd}(\Gamma) = \mathbb{Q}.$$ 

For the underlying system we let the set $A$ consist of two symbols, say

$$A := \{a, b\}$$

with $a \neq b$. The intervals are

$$I_a := (0, 1) \quad \text{and} \quad I_b := (1, \infty).$$

The sets of parabolic elements are

$$P_{a,a} := P_{b,b} := \emptyset, \quad P_{a,b} := \{p_1\}, \quad P_{b,a} := \{p_2\}$$

with

$$p_1 := \begin{bmatrix} 1 & -1 \\ 0 & 1 \end{bmatrix}, \quad p_2 := \begin{bmatrix} 1 & 0 \\ -1 & 1 \end{bmatrix}.$$ 

For all $p \in P_{a,b} \cup P_{b,a}$ we set $g_p := \text{id}$. For all $n \in \mathbb{N}$ we have

$$p_1^{-n} \cdot I_{a,\text{st}} = (n, n + 1)_{\text{st}} \subseteq I_{b,\text{st}},$$

$$p_2^{-n} \cdot I_{b,\text{st}} = \left(1 + \frac{1}{n}, \frac{1}{n}\right)_{\text{st}} \subseteq I_{a,\text{st}}.$$ 

We refer to [41] for the proof that Properties 1–5 are satisfied.

The induced discrete dynamical system is

$$p_1^{-n} \cdot I_{a,\text{st}} \times \{b\} \to I_{a,\text{st}} \times \{a\},$$

$$(x, b) \mapsto (p_1^n \cdot x, a) = (x - n, a),$$
meromorphic continuation of Selberg zeta functions with... Page 25 of 55

\[ p_2^{-n} I_b, st \times \{a\} \rightarrow I_b, st \times \{b\}, \]

\[ (x, a) \mapsto (p_2^n x, b) = \left( \frac{x}{-nx + 1}, b \right). \]

The transfer operator is

\[
\mathcal{L}_s = \begin{pmatrix}
0 & \sum_{n \in \mathbb{N}} \alpha_s(p_1^n) \\
\sum_{n \in \mathbb{N}} \alpha_s(p_2^n) & 0
\end{pmatrix}
\]  (30)

acting on functions (with sufficient regularity) of the form

\[ f = \left( f_a : I_{a, st} \rightarrow V, f_b : I_{b, st} \rightarrow V \right). \]

The relation of the transfer operator \( \mathcal{L}_s \) to Mayer’s transfer operator \( \mathcal{L}_s^M \) is given by diagonalization, and taking the representation \( \chi \) to be trivial and one-dimensional. Changing the basis of the function space shows that \( \mathcal{L}_s \) is conjugate to

\[
\left( \begin{array}{cc}
\mathcal{L}_s^M & 0 \\
0 & -\mathcal{L}_s^M
\end{array} \right).
\]

Geometrically, this conjugation means that the geodesic flow on \( \text{PSL}_2(\mathbb{Z}) \setminus \mathbb{H} \) is split into two copies of the billiard flow on \( \text{PGL}_2(\mathbb{Z}) \setminus \mathbb{H} \), one with Neumann boundary value conditions and one with Dirichlet boundary value conditions. We refer to [49] for a more detailed discussion.

4.4 Statement of Main Theorem for meromorphic continuability (strong version)

With the preparations from Sects. 4.1–4.3 we are now able to state the strong version of the main theorem regarding the meromorphic continuability of twisted Selberg zeta functions.

For any Banach space \( \mathcal{B} \), we let

\[ \text{Op}(\mathcal{B}) := \{ L : \mathcal{B} \rightarrow \mathcal{B} \text{ linear, bounded} \} \]

denote the bounded linear operators on \( \mathcal{B} \).

**Theorem 4.2** Let \( \Gamma \) be a geometrically finite Fuchsian group which admits a strict transfer operator approach, and let \( \chi : \Gamma \rightarrow \text{GL}(V) \) be a finite-dimensional representation of \( \Gamma \) with non-expanding cusp monodromy. Let

\[ \mathcal{S} := (A, (I_a)_{a \in A}, (P_{a,b})_{a,b \in A}, (C_{a,b})_{a,b \in A}, ((g_p)_{p \in P_{a,b}})_{a,b \in A}) \]

be a structure tuple as in (22), and let \( \mathcal{B}(I; V) \) denote the associated Banach space as in (29). For \( s \in \mathbb{C} \) let \( \mathcal{L}_s \) denote the associated (formal) transfer operator. Further,
for \( p \in \Gamma \) parabolic let \( \text{jc}(p) \) denote the family (with multiplicities) of the lengths of Jordan chains for eigenvalue 1 of \( \chi(p) \), and let

\[
d_1 := \max\{\text{jc}(p) \mid p \in \Gamma \text{ parabolic}\}
\]

denote the maximal length of a Jordan chain ranging over all \( \chi(p) \), \( p \in \Gamma \) parabolic, with eigenvalue 1. Further, let \( d_0 \) denote the maximal length of a Jordan chain ranging over all \( \chi(p) \), \( p \in \Gamma \) parabolic, and all eigenvalues.

Then the following properties hold:

(i) For \( \Re s > \frac{d_0}{2} \), the (formal) transfer operator \( L_s \) defines an (actual) operator on \( \mathcal{B}(I; V) \). It is bounded and nuclear of order 0.

(ii) The map

\[
\left\{ s \in \mathbb{C} \mid \Re s > \frac{d_0}{2} \right\} \to \text{Op}(\mathcal{B}(I; V)), \quad s \mapsto L_s
\]

extends meromorphically to all of \( \mathbb{C} \) with values in nuclear operators of order 0 in \( \mathcal{B}(I; V) \). All poles are simple and of the form \( \frac{1}{2}(d_1 - k) \) with \( k \in \mathbb{N}_0 \). For each pole \( s_0 \), there exists a neighborhood \( U \) of \( s_0 \) such that the meromorphic continuation \( \tilde{L}_s \) is of the form

\[
\tilde{L}_s = \frac{1}{s - s_0} A_s + B_s,
\]

where the operators \( A_s \) and \( B_s \) are holomorphic on \( U \), and \( A_s \) is of rank at most

\[
\sum_{a, b \in A} \sum_{p \in P_{a, b}} \sum_{d \in \text{jc}(p)} \left( \frac{d + 1}{2} \right).
\]

(iii) For \( \Re s \gg 1 \), \( Z_{\Gamma, \chi}(s) = \det(1 - L_s) \).

(iv) The Selberg zeta function \( Z_{\Gamma, \chi} \) extends meromorphically to all of \( \mathbb{C} \). Its poles are contained in \( \frac{1}{2}(d_1 - \mathbb{N}_0) \). The order of any pole is at most

\[
\sum_{a, b \in A} \sum_{p \in P_{a, b}} \sum_{d \in \text{jc}(p)} \left( \frac{d + 1}{2} \right).
\]

For unitary representations \( \chi \), the proof of Theorem 4.2 is by now standard \[16,38,50,52\]. Thus, our main emphasis is on discussing the necessary changes and extensions of the standard proof in order to accommodate arbitrary representations with non-expanding cusp monodromy.

The proof of Theorem 4.2 is split into several parts: In Sect. 4.5 below we will show Theorem 4.2(i). Compared to the standard proof the essential part to show for Theorem 4.2(i) is the convergence of the infinite sums featuring in the transfer operator, see (28). Every such infinite sum is governed by a parabolic element. Hence, \textit{a priori},
the growth of $\chi(p^n)$ as $n \to \infty$ might cause divergence. However, the logarithm-type law Lemma 3.3 gives us sufficient control of this growth in order to establish convergence.

In Sect. 4.6 below we will prove Theorem 4.2(iii). The key result for this step is an easy-to-evaluate formula for the trace of $\alpha_s(h)$, $h \in \Gamma$ hyperbolic, which we will provide in Lemma 4.4 below.

In Sect. 4.7 below we will show Theorem 4.2(ii). This result is based on a reduction to the Lerch transcendent and its derivatives, which allows us to handle the contribution of non-diagonalizable families of endomorphisms $(\chi(p^n))_{n \in \mathbb{N}}$ for $p \in \Gamma$ parabolic.

We remark that for some Fuchsian groups, in particular for Fuchsian Schottky groups, the standard setup uses Hilbert spaces of holomorphic $L^2$-functions in place of $\mathcal{B}(E_A; V)$ [22]. The proof of Theorem 4.2 applies, *mutatis mutandis*, to this situation. We will not discuss it separately.

Throughout Sects. 4.5–4.8 we will use, without any further reference, the notation from Sects. 4.1–4.4, in particular regarding the structure tuple $S$ and the transfer operators $\mathcal{L}_s$. Further we fix a family $E_A = (E_a)_{a \in A}$ as provided by Property 5.

### 4.5 Nuclearity of transfer operators

In Proposition 4.3 below we will show that $\mathcal{L}_s$ defines a nuclear operator of order 0 on $\mathcal{B}(E_A; V)$, which then proves Theorem 4.2(i).

Essential for the proof of Proposition 4.3 is to understand the influence of the families $(\chi(p^n))_{n \in \mathbb{N}}$, $p \in \Gamma$ parabolic, on the growth of the transfer operators $\mathcal{L}_s$. For this, we make crucial use of Lemma 3.3.

**Proposition 4.3** Let $s \in \mathbb{C}$, $\text{Re } s > \frac{d_0}{2}$. Then the transfer operator $\mathcal{L}_s$ defines a bounded self-map of $\mathcal{B}(E_A; V)$. As such $\mathcal{L}_s$ is nuclear of order 0.

**Proof** For any open subset $\mathcal{D} \subseteq \mathbb{C}$ we let

$$\mathcal{H}(\mathcal{D}; V) := \{ f : \mathcal{D} \to V \text{ holomorphic} \}$$

denote the space of holomorphic functions $\mathcal{D} \to V$ and endow it with the compact-open topology, which turns $\mathcal{H}(\mathcal{D}; V)$ into a nuclear space [19, I, §2, no. 3, Corollary]. We further set

$$\mathcal{H}(E_A; V) := \bigoplus_{a \in A} \mathcal{H}(E_a; V),$$

and endow it with the product space structure. Hence, $\mathcal{H}(E_A; V)$ is a nuclear space.

We start by showing that the transfer operator $\mathcal{L}_s$ defines a map from $\mathcal{H}(E_A; V)$ to $\mathcal{B}(E_A; V)$. To that end let

$$f = \bigoplus_{a \in A} f_a \in \mathcal{H}(E_A; V)$$
and set

\[ h := \mathcal{L}_s f, \quad h = \bigoplus_{a \in A} h_a. \]

Let \( a \in A \). Then

\[ h_a = \sum_{b \in A} \mathcal{L}_{s,a,b} f_b \]

\[ = \sum_{b \in A} \left( \sum_{g \in C_{a,b}} \alpha_s(g) f_b + \sum_{p \in P_{a,b}, n \in \mathbb{N}} \alpha_s(g_p p^n) f_b \right). \quad (31) \]

We claim that \( h_a \in \mathcal{B}(E_a; V) \). Let \( b \in A \). Property 5(ii)-(iii) immediately imply that for all \( g \in C_{a,b} \), the map

\[ z \mapsto \alpha_s(g) f_b(z) = \left( (g^{-1}')(z) \right)^s \chi(g) f_b(g^{-1}z) \quad (32) \]

is well-defined on \( \overline{E}_a \), continuous on \( \overline{E}_a \) and holomorphic on \( E_a \). This and the finiteness of \( C_{a,b} \) implies that the first summand in (31) defines an element of \( \mathcal{B}(E_a; V) \). Therefore it remains to show that for each \( p \in P_{a,b} \), the infinite sum

\[ F := F_{s,p,b} := \sum_{n \in \mathbb{N}} \alpha_s(g_p p^n) f_b \]

defines an element of \( \mathcal{B}(E_a; V) \). Recall that \( d_0 \) denotes the maximal length of a Jordan chain of \( \chi(q) \), where \( q \) ranges over all parabolic elements in \( \Gamma \). Lemma 3.3 shows

\[ \| \chi(p^n) \| \ll n^{d_0-1} \]

for all \( n \in \mathbb{N} \). Further, by Property 5(iv), there exists a compact subset \( K_{a,b,p} \) of \( \hat{\mathbb{C}} \) such that for each \( n \in \mathbb{N} \),

\[ p^{-n} g_p^{-1} E_a \subseteq K_{a,b,p} \subseteq E_b. \quad (33) \]

Since \( f_b \) is continuous on \( E_b \), \( \| f_b(p^{-n} g_p^{-1} z) \| \) is bounded on \( \overline{E}_a \), uniformly in \( n \in \mathbb{N} \). Thus,

\[ \sup \left\{ \| \chi(g_p p^n) f_b(p^{-n} g_p^{-1} z) \| : z \in \overline{E}_a, \ n \in \mathbb{N} \right\} \ll n^{d_0-1}. \]

Since \( p \) is parabolic, and hence conjugate to \( \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} \) in \( \text{PSL}_2(\mathbb{R}) \), there exist \( c, d \in \mathbb{R} \) such that

\[ p^n = \begin{bmatrix} 1 + ncd & nd^2 \\ -nc^2 & 1 - ncd \end{bmatrix} \quad (34) \]

for all \( n \in \mathbb{Z} \). The boundedness of \( K_{a,b,p} \) in (33) implies that \( c \neq 0 \).
Let

$$g_p = \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix}. $$

Then

$$p^{-n}g_p^{-1} = \begin{bmatrix} * & * \\ g_{22}nc^2 - g_{21}(1 + ncd) & -g_{12}nc^2 + g_{11}(1 + ncd) \end{bmatrix}. $$

By Property 5(iv) or, equivalently, by (33),

$$p^{-n}g_p^{-1}.z \neq \infty$$

for all \(z \in \overline{E}_a, n \in \mathbb{N},\) and hence

$$\text{Re} \left((g_{22}nc^2 - g_{21}(1 + ncd))z - g_{12}nc^2 + g_{11}(1 + ncd)\right) \neq 0$$

Thus, for all \(z \in \overline{E}_a, n \in \mathbb{N},\)

$$\left((p^{-n}g_p^{-1})'(z)\right)^s = \left((nc\left(c(g_{22}z - g_{12}) + d(-g_{21}z + g_{11})\right) - g_{21}z + g_{11})^{-2}\right)^s$$

is well-defined, and

$$\left|\left((p^{-n}g_p^{-1})'(z)\right)^s\right| \ll_s |nc\left(c(g_{22}z - g_{12}) + d(-g_{21}z + g_{11})\right) - g_{21}z + g_{11}|^{-2\text{Re}s}. $$

By Property 5(v),

$$c(g_{22}z - g_{12}) + d(-g_{21}z + g_{11}) \neq 0$$

for all \(z \in \overline{E}_a.\) In combination with \(c \neq 0\) (see above) it follows that

$$\sup_{z \in \overline{E}_a} \left|\left((p^{-n}g_p^{-1})'(z)\right)^s\right| \ll_s n^{-2\text{Re}s}. $$

Therefore

$$\sup_{z \in \overline{E}_a} \|\alpha_s(g_pp^n)f_b(z)\| \ll_s n^{d_0 - 1 - 2\text{Re}s}. \quad (35)$$

This implies that if \(\text{Re}\, s > \frac{d_0}{2},\) the sum \(F = \sum_{n \in \mathbb{N}} \alpha_s(g_pp^n)f_b\) converges uniformly on \(\overline{E}_a.\) By Weierstrass' Theorem, \(F\) is continuous on \(\overline{E}_a\) and holomorphic on \(E_a.\) Therefore \(F \in \mathcal{B}(\mathcal{E}_a; V),\) and hence \(\mathcal{L}_s : \mathcal{H}(\mathcal{E}_a; V) \to \mathcal{B}(\mathcal{E}_a; V).\)

As in the standard proof [16,38,50,52], one can now deduce that, as a map from \(\mathcal{H}(\mathcal{E}_a; V)\) to \(\mathcal{B}(\mathcal{E}_a; V),\) \(\mathcal{L}_s\) is nuclear of order 0. Then one further concludes that \(\mathcal{L}_s\) remains nuclear of order 0 when considered as a map \(\mathcal{B}(\mathcal{E}_a; V) \to \mathcal{B}(\mathcal{E}_a; V).\) \(\square\)
4.6 The twisted Selberg zeta function as a Fredholm determinant

In this section we will show that, for $\Re s \gg 1$, the twisted Selberg zeta function $Z_{\Gamma, \chi}$ equals the Fredholm determinant of the transfer operator family $L_s$. This immediately proves Theorem 4.2(iii).

In order to simplify the proof of the trace formula in Lemma 4.4 below, we set

$$\tau_s^C(h^{-1})\varphi(x) := (h'(x))^s \varphi(h \cdot x)$$

for subsets $U \subseteq \mathbb{R}$, functions $\varphi : U \to \mathbb{C}$, $h \in \Gamma$ and $x \in U$, whenever it makes sense.

**Lemma 4.4** Let $\mathcal{D}$ be an open, connected, simply connected, bounded subset of $\hat{\mathbb{C}}$ and suppose that $h \in \Gamma$ satisfies $h^{-1}.\mathcal{D} \subseteq \mathcal{D}$. Then the map

$$\alpha_s(h) : B(\mathcal{D}; V) \to B(\mathcal{D}; V)$$

is nuclear of order 0 with trace

$$\text{Tr} \alpha_s(h) = \frac{N(h)^{-s}}{1 - N(h)^{-1}} \text{tr} \chi(h).$$

**Proof** Note that the element $h$ necessarily is hyperbolic. Thus, the norm $N(h)$ is indeed defined.

If $(V, \chi)$ is the trivial one-dimensional representation, then $\alpha_s(h) = \tau_s^C(h)$. Note that $N(h) = N(h^{-1})$. By [52] we have

$$\text{Tr} \tau_s^C(h) = \frac{N(h)^{-s}}{1 - N(h)^{-1}},$$

which shows the statement of the Lemma for this particular case.

Let $(V, \chi)$ now be a generic representation with non-expanding cusp monodromy and fix a basis $B_V := \{v_1, \ldots, v_d\}$ of $V$. The map

$$B(\mathcal{D}; \bigotimes^d V) \to B(\mathcal{D}; V), \quad (f_j)_{j=1}^d \mapsto \sum_{j=1}^d f_j v_j$$

is an isomorphism of Banach spaces. Let

$$M = (m_{ij})_{i,j=1,\ldots,d}$$

be the $d \times d$-matrix that represents $\chi(h)$ with respect to $B_V$. Under the isomorphism (38), the action of $\alpha_s(h)$ is then given by

$$\left( m_{ij} \tau_s^C(h) \right)_{i,j=1,\ldots,d}.$$
Since

\[ B(D; \mathbb{C})^d \cong \bigoplus_{j=1}^d B(D; \mathbb{C})v_j \]

is a topological direct sum, [20, Chapitre II.3, Prop. 2] shows that

\[ \text{Tr}_\alpha_s(h) = \sum_{j=1}^d m_{jj} \text{Tr}_{\tau_s^j}^C(h). \]

Combining this equality with (37) completes the proof. \(\square\)

**Proposition 4.5** For \( \text{Re} \, s \gg 1 \), 

\[ Z_{\Gamma, \chi}(s) = \det(1 - L_s). \]

**Proof** Recall the constant \( C = C(\Gamma, \chi) \) from Theorem 3.1, and that \( d_0 \) denotes the maximal length of a Jordan block of \( \chi(p) \), \( p \in \Gamma \) parabolic. Let \( s \in \mathbb{C} \) with \( \text{Re} \, s > \max\{C, d_0/2\} \). By Theorem 3.1, \( Z_{\Gamma, \chi}(s) \) converges. Further Proposition 4.3 implies that the Fredholm determinant of \( L_s \) exists. Analogously to the case of the trivial one-dimensional representation one sees that

\[ \det(1 - L_s) = \exp \left( -\sum_{n \in \mathbb{N}} \frac{1}{n} \text{Tr} L_s^n \right). \]

Property 5 combined with (26) shows that for \( n \in \mathbb{N}, h \in \text{Per}_n \) and \( a \in A \) as in (26) we have

\[ h^{-1} E_a \subseteq E_a. \]

By [20, Chapitre II.3, Prop. 2] and Lemma 4.4,

\[ \text{Tr} L_s^n = \sum_{h \in \text{Per}_n} \text{Tr} \alpha_s(h) = \sum_{h \in \text{Per}_n} \frac{N(h)^{-s}}{1 - N(h)^{-1}} \text{tr} \chi(h). \]

Then Properties 3 and 4 imply

\[ \log Z_{\Gamma, \chi}(s) = \sum_{[g] \in [\Gamma]_p} \sum_{k=0}^\infty \log \left( \det \left( 1 - \chi(g)N(g)^{-(s+k)} \right) \right) \]

\[ = - \sum_{[g] \in [\Gamma]_p} \sum_{m=1}^\infty \frac{1}{m} \frac{N(g^m)^{-s}}{1 - N(g^m)^{-1}} \text{tr} \chi(g^m) \]

\[ = - \sum_{[g] \in [\Gamma]_h} \frac{1}{m(g)} \frac{N(g)^{-s}}{1 - N(g)^{-1}} \text{tr} \chi(g), \]
\[
\begin{align*}
&= - \sum_{[g] \in [\Gamma]_h} p(g) \frac{N(g)^{-s}}{w(g)} \frac{1}{1 - N(g)^{-1}} \text{tr} \chi(g) \\
&= - \sum_{w=1}^{\infty} \frac{1}{w} \sum_{h \in \text{Per}_w} \frac{N(h)^{-s}}{1 - N(h)^{-1}} \text{tr} \chi(h) \\
&= - \sum_{w=1}^{\infty} \frac{1}{w} \text{Tr} \mathcal{L}_s^w \\
&= \log \left( \det \left( 1 - \mathcal{L}_s \right) \right).
\end{align*}
\]

This completes the proof. \qed

### 4.7 Meromorphic continuation of the transfer operator family

In this section we will show that the map

\[
\left\{ s \in \mathbb{C} \ \bigg| \ \text{Re} \ s > \frac{d_0}{2} \right\} \to \text{Op} \left( \mathcal{B}(\mathcal{E}_A; V) \right), \quad s \mapsto \mathcal{L}_s,
\]

extends meromorphically to all of \( \mathbb{C} \) with values in nuclear operators of order 0 on \( \mathcal{B}(\mathcal{E}_A; V) \). This means that there exists a discrete set \( P \subseteq \mathbb{C} \) (‘candidates for poles’) and for each \( s \in \mathbb{C} \setminus P \) we find an operator \( \tilde{\mathcal{L}}_s : \mathcal{B}(\mathcal{E}_A; V) \to \mathcal{B}(\mathcal{E}_A; V) \) which is nuclear of order 0 and which equals \( \mathcal{L}_s \) for \( \text{Re} \ s > \frac{d_0}{2} \). Moreover, for each \( f \in \mathcal{B}(\mathcal{E}_A; V) \) and each point

\[
z \in \mathcal{E} := \bigcup_{a \in A} \mathcal{E}_a
\]

the map \( s \mapsto \tilde{\mathcal{L}}_s f(z) \) is meromorphic with poles contained in \( P \), and the map \( (s, z) \mapsto \tilde{\mathcal{L}}_s f(z) \) is continuous on \( (\mathbb{C} \setminus P) \times \mathcal{E} \).

If \( \chi \) is the trivial one-dimensional representation, then the Hurwitz zeta function plays a crucial role in the proof of this meromorphic continuability. For unitary representations \( \chi \), the Lerch zeta function features. For generic representations \( \chi \) with non-expanding cusp monodromy, we will take advantage of the Lerch transcendent and its derivatives. The main emphasis of this section is to show how the Lerch transcendent allows us to generalize the standard proof for trivial and unitary representations to the larger class of representations with non-expanding cusp monodromy. This result shows Theorem 4.2(ii).

We denote the Lerch transcendent by

\[
\Phi(s, \lambda, w) := \sum_{n=0}^{\infty} \frac{\lambda^n}{(n + w)^s},
\]
For \( m \in \mathbb{N}_0 \) we set
\[
\Phi_m(s, \lambda, w) := \frac{1}{m!} \partial^m_\lambda \Phi(s, \lambda, w) = \sum_{n=m}^{\infty} \binom{n}{m} \frac{\lambda^{n-m}}{(n+w)^s}.
\]
For \(|\lambda| = 1\) and \( w \notin -\mathbb{N}_0 \) (the values we need) the series converges and defines a holomorphic function on \( \{ s \in \mathbb{C} \mid \Re s > m + 1 \} \). For \( \lambda \neq 1 \) it extends in \( s \) to an entire function. For \( \lambda = 1 \) it extends meromorphically in \( s \) to all of \( \mathbb{C} \) with all poles contained in \( \{ 1, \ldots, m + 1 \} \), all simple, and \( m + 1 \) being a pole. In both cases, the meromorphic extension is continuous on its domain of holomorphy as a function in the two variables \((s, w)\). To be more precise, the extension is unique for \( w \in \mathbb{C} \setminus (-\infty, 0] \).

The additional extension in the \( w \)-variable requires a choice of an extension of the complex logarithm.

For \( \Phi = \Phi_0 \) proofs of these statements can be found in the literature, for \( \Phi_m \) with \( m > 0 \) they can be deduced from the one for \( \Phi_0 \) or shown in analogy. For the convenience of the reader, we provide a proof in Sect. 1 below. We refer to [35] for detailed discussions on the meromorphic continuability of the Lerch transcendent as a function of all three variables.

**Proposition 4.6** The map \( s \mapsto L_s \) extends to a meromorphic function on \( \mathbb{C} \) with values in nuclear operators of order 0 on \( B(\mathcal{E}_A; V) \). The possible poles are all simple and contained in \( \frac{1}{2}(d_1 - \mathbb{N}_0) \) (recall \( d_1 \) from Theorem 4.2). For each pole \( s_0 \), there is a neighborhood \( U \) of \( s_0 \) such that the meromorphic continuation \( \tilde{L}_s \) is of the form
\[
\tilde{L}_s = \frac{1}{s - s_0} A_s + B_s,
\]
where the operators \( A_s \) and \( B_s \) are holomorphic on \( U \), and \( A_s \) is of rank at most
\[
\sum_{a, b \in A} \sum_{p \in P_{a, b}} \sum_{d \in \mathcal{J}(p)} \left( \frac{d + 1}{2} \right).
\]

**Proof** For Banach spaces \( B_1, B_2 \) we denote by \( \text{Op}(B_1; B_2) \) the space of all linear operators \( B_1 \rightarrow B_2 \).

It suffices to show that for \( a, b \in A \) and \( p \in P_{a, b} \) the map
\[
\{ s \in \mathbb{C} \mid \Re s > d_0/2 \} \rightarrow \text{Op} \left( B(\mathcal{E}_b; V); B(\mathcal{E}_a; V) \right), \quad s \mapsto L_s := \sum_{n \in \mathbb{N}} \alpha_s(g_p p^n)
\]
extends meromorphically to all of \( \mathbb{C} \) with values in nuclear operators of order 0 and all poles contained in \( P \). Since
\[
\sum_{n \in \mathbb{N}} \alpha_s(g_p p^n) = \alpha_s(g_p) \sum_{n \in \mathbb{N}} \alpha_s(p^n),
\]
...
it also suffices to show that
\[
\{ s \in \mathbb{C} \mid \text{Re } s > \frac{d_0}{2} \} \rightarrow \text{Op} \left( \mathcal{B}(\mathcal{E}_b; V); \alpha_s(g_p^{-1})\mathcal{B}(\mathcal{E}_a; V) \right), \quad s \mapsto \sum_{n \in \mathbb{N}} \alpha_s(p^n)
\]
extends meromorphically to all of \( \mathbb{C} \). Hence, without loss of generality, we may assume that \( g_p = \text{id} \).

Let \( a, b \in A \) and \( p \in P_{a,b} \). Represent \( \chi(p) \) in Jordan normal form and fix a Jordan block. Let \( \lambda \in \mathbb{C} \) be the eigenvalue of this Jordan block, and let \( d_b \) denote the size of the Jordan block (that is, the length of the defining Jordan chain). It suffices to study the meromorphic continuability of the maps
\[
\{ s \in \mathbb{C} \mid \text{Re } s > \frac{d_0}{2} \} \rightarrow \text{Op} \left( \mathcal{B}(\mathcal{E}_b; \mathbb{C}); \mathcal{B}(\mathcal{E}_a; \mathbb{C}) \right)
\]
\[
s \mapsto H_{s,m} := \sum_{n \in \mathbb{N}} \binom{n}{m} \lambda^{n-m} r_s(c)(p^n)
\]
for \( m \in \{0, 1, \ldots, d_b - 1\} \). Recall that by hypothesis we have \( |\lambda| = 1 \).

Recall that (without loss of generality) \( \mathcal{E}_a \) and \( \mathcal{E}_b \) are subsets of \( \mathbb{C} \). For \( M \in \mathbb{N}_0 \), let
\[
Q_M, P_M : \mathcal{B}(\mathcal{E}_b; \mathbb{C}) \rightarrow \mathcal{B}(\mathcal{E}_b; \mathbb{C})
\]
be the operator that maps a function to its Taylor polynomial of degree \( M \) centered at the fixed point \( x_p \) of \( p \) or subtracts this Taylor polynomial from it, respectively. Thus we set
\[
Q_M(f)(z) := \sum_{k=0}^{M} \frac{f^{(k)}(x_p)}{k!} (z - x_p)^k
\]
and
\[
P_M(f)(z) := (1 - Q_M)f(z) := f(z) - \sum_{k=0}^{M} \frac{f^{(k)}(x_p)}{k!} (z - x_p)^k.
\]
Then
\[
H_{s,m} = H_{s,m} \circ Q_M + H_{s,m} \circ P_M.
\]
We first study \( H_{s,m} \circ Q_M \). By Property 5(iv), \( \overline{\mathcal{E}_a} \) does not contain the fixed point \( x_p \) of \( p \). Using the explicit form (34) for \( p^n, n \in \mathbb{N}_0 \), we find \( x_p = -d/c \). Thus, for \( z \in \mathcal{E}_a \) we have \( cz + d \neq 0 \) and
\[
p^{-n}z - x_p = \left( c^2 \left( n + \frac{1}{c(cz + d)} \right) \right)^{-1}.
\]
By Property 5(iv),

\[ \text{Re} \left( nc(cz + d) + 1 \right) \neq 0 \]

for all \( z \in \mathcal{E}_a, n \in \mathbb{N} \). Thus,

\[
\left( (p^{-n})' (z) \right)^s = \left( (nc(cz + d) + 1)^{-2} \right)^s \left( n + \frac{1}{c(cz + d)} \right)^{-2s}
\]

for a function \( \varphi : \mathbb{N} \times \mathcal{E}_a \rightarrow 2\pi \mathbb{Z} \). The function \( \varphi \) is continuous and hence constant in the \( z \)-variable. Further, since the argument (or phase) of \( n + \frac{1}{c(cz + d)} \) tends to 0 as \( n \rightarrow \infty \), there exists \( n_0 \in \mathbb{N} \) such that \( \varphi \) is constant on \( \mathbb{N} \geq n_0 \times \mathcal{E}_a \) (note that \( \frac{1}{c(cz + d)} \) is bounded on \( \mathcal{E}_a \)). Without loss of generality we may assume that \( \varphi = 0 \), and that \( \text{Re} \left( c(cz + d) \right) > 0 \) for all \( z \in \mathcal{E}_a \). Hence,

\[
\left( (p^{-n})' (z) \right)^s = (c(cz + d))^{-2s} \left( n + \frac{1}{c(cz + d)} \right)^{-2s}.
\]

Then

\[
(H_{s,m} \circ Q_M) f(z) = \sum_{k=0}^{M} \frac{f^{(k)}(x_p)}{k!} \sum_{n \in \mathbb{N}} \binom{n}{m} \lambda^{-m} \left( (p^{-n})' (z) \right)^s (p^{-n} \cdot z - x_p)^k
\]

\[
= \sum_{k=0}^{M} \frac{f^{(k)}(x_p)}{k!} c^{-2} (c(cz + d))^{-2s} \sum_{n \in \mathbb{N}} \binom{n}{m} \lambda^{-m} \left( n + \frac{1}{c(cz + d)} \right)^{-2s-k}
\]

\[
= \sum_{k=0}^{M} \frac{f^{(k)}(x_p)}{k!} c^{-2} (c(cz + d))^{-2s} \Phi_m \left( 2s + k, \lambda, \delta_{0,m} + \frac{1}{c(cz + d)} \right),
\]

where

\[
\delta_{0,m} := \begin{cases} 
1 & \text{if } m = 0 \\
0 & \text{otherwise.}
\end{cases}
\]

The properties of \( \Phi_m \), as recalled just before the statement of this proposition, yield the existence of a meromorphic continuation of \( s \mapsto (H_{s,m} \circ Q_M) f(z) \) with the localization and order of poles as claimed. Note that \( \Phi_m \) may only contribute a pole if \( \lambda = 1 \), and hence the first possible pole is located at \( s = d_1/2 \) instead of \( d_0/2 \).

The extension of \( H_{s,m} \circ Q_M \) is nuclear of order 0 as a finite rank operator.
We now study $H_{s,m} \circ P_M$. We fix an open ball $B$ in $\mathbb{C}$ around $x_p$ such that $B \subseteq \mathcal{E}_b$. Let $f \in \mathcal{B}(\mathcal{E}_b)$ and set $f_M := P_M(f)$. For all $z \in B$ we have

$$|f_M(z)| \ll |z - x_p|^{M+1}.$$  

Further, for all $z \in \mathbb{C}$, we have

$$\lim_{n \to \infty} p^{-n}z = x_p.$$  

Let $z \in \mathcal{E}_a$. Then there exists $n_0 = n_0(z) \in \mathbb{N}$ such that

$$p^{-n}z \in B$$  

for $n > n_0$. Thus,

$$\left|(H_{s,m} \circ P_M) f(z)\right| = \left|\sum_{n \in \mathbb{N}} \binom{n}{m} \lambda^{n-m} \tau_s^C(p^n) f_M(z)\right|$$

$$\leq \left|\sum_{n \leq n_0} \binom{n}{m} \lambda^{n-m} \tau_s^C(p^n) f_M(z)\right|$$

$$+ \sum_{n > n_0} \binom{n}{m} \left|(p^{-n})'(z)\right|^{\Re s} \left|f_M(p^{-n}z)\right|$$

$$\ll \left|\sum_{n \leq n_0} \binom{n}{m} \lambda^{n-m} \tau_s^C(p^n) f_M(z)\right|$$

$$+ \sum_{n > n_0} n^{d_0 - 1} |cz + d|^{-2 \Re s} \cdot n + \frac{1}{c(cz + d)}^{-(M+1+2 \Re s)}.$$  

Hence, $(H_{s,m} \circ P_M) f(z)$ converges for

$$\Re s > \frac{d_0 - M - 1}{2}.$$  

The operator $H_{s,m} \circ P_M$ is nuclear of order 0 since $P_M$ is bounded.

Since these arguments apply to each $M \in \mathbb{N}_0$, the proof is complete.\qed

4.8 Proof of Theorem 4.2

Theorem 4.2(i) is just Proposition 4.3, and Theorem 4.2(iii) is shown by Proposition 4.5. Theorem 4.2(ii) follows from Proposition 4.6. The remaining statement Theorem 4.2(iv) are implied by (i)–(iii) analogously to the case of trivial or unitary representations (see, e.g., the proof of [41, Theorem 4.14] or of [50, Theorem 5.1]).\qed
5 Representations with non-expanding cusp monodromy

Obviously, there are (non-uniform) Fuchsian groups $\Gamma$ that admit finite-dimensional representations without non-expanding cusp monodromy. In this section we will provide examples of finite-dimensional representations with non-expanding cusp monodromy.

For Fuchsian groups $\Gamma$ without any parabolic elements (in particular, for uniform Fuchsian groups), the condition on parabolic elements is void and hence any representation of $\Gamma$ has non-expanding cusp monodromy. For non-uniform Fuchsian groups $\Gamma$, obvious examples of representations with non-expanding cusp monodromy are all unitary representations, or more generally, all representations that are unitary in the cusps. We provide more examples, and in particular also non-obvious ones, in what follows.

With Proposition 5.1 below we will show that every representation of $\Gamma$ that is the restriction of a finite-dimensional representation of $\text{SL}_2(\mathbb{R})$ has non-expanding cusp monodromy. This result is complemented by Remark 5.2 below that proves that this construction indeed yields infinitely many representations of $\Gamma$.

In Proposition 5.3 below we will show that for every geometrically finite Fuchsian group $\Gamma$ with at least one parabolic element, the class of finite-dimensional representations with non-expanding cusp monodromy contains more elements than provided by Proposition 5.1 and the obvious examples from above. More precisely, we will construct a finite-dimensional representation of $\Gamma$ with non-expanding cusp monodromy that is neither unitary in the cusps nor a representation from Proposition 5.1 nor a direct sum of representations of smaller dimensions.

Finally, Proposition 5.4 below will show for all subgroups $\Lambda$ of $\Gamma$ of finite index that all finite-dimensional representations with non-expanding cusp monodromy of $\Lambda$ induce to finite-dimensional representations with non-expanding cusp monodromy of $\Gamma$.

**Proposition 5.1** Let $\chi : \text{SL}_2(\mathbb{R}) \to \text{GL}(V)$ be a finite-dimensional representation of $\text{SL}_2(\mathbb{R})$ and let $\Gamma$ be a Fuchsian group. If $\chi$ descends to a representation of $\text{PSL}_2(\mathbb{R})$ then the restriction $\chi|_\Gamma$ of $\chi$ to $\Gamma$ has non-expanding cusp monodromy.

**Proof** Since every finite-dimensional representation of $\text{SL}_2(\mathbb{R})$ decomposes into a direct sum of irreducible representations (see [24, Corollary 2.2]), it suffices to prove the statement for irreducible representations. As it is well known (see, e.g., [24, (2.1), Proposition 2.1 (a) and (c), Remark after Proposition 2.1 and the first line in the proof of Proposition 2.1]), up to equivalence, the irreducible finite-dimensional representations of $\text{SL}_2(\mathbb{R})$ are parametrized by non-negative integers. The representation $(V_n, \chi_n)$ with parameter $n \in \mathbb{N}_0$ can be realized on the complex vector space $V_n$ of the polynomials in two complex variables $x$ and $y$ that are homogeneous of degree $n$. The action is then given by

$$\chi_n(g)(f)(\begin{pmatrix} x \\ y \end{pmatrix}) = f(g^{-1}(\begin{pmatrix} x \\ y \end{pmatrix})), \quad (39)$$

where $g \in \text{SL}_2(\mathbb{R})$ and $f \in V_n$.

In order to show that $\chi_n$ has non-expanding cusp monodromy, it is sufficient to determine the eigenvalues of $\chi(p)$ for the parabolic element $p = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$. A basis for
$V_n$ is given by the $n + 1$ monomials $e_0, \ldots, e_n$ of degree $n$ in 2 variables:

$$e_j \left( \frac{x}{y} \right) := x^j y^{n-j}, \quad j = 0, \ldots, n.$$ 

A straightforward calculation shows

$$\chi_n(p) e_j = e_j + \sum_{k=1}^{j} \binom{j}{k} e_{j-k} \quad (40)$$

for all $j \in \{0, 1, \ldots, n\}$, which immediately implies (e.g., by representing $\chi_n(p)$ with respect to the basis $\{e_0, \ldots, e_n\}$) that 1 is the only eigenvalue of $\chi_n(p)$. Hence $\chi_n$ has non-expanding cusp monodromy.

**Remark 5.2** Proposition 5.1 indeed yields finite-dimensional representations with non-expanding cusp monodromy of Fuchsian groups. As soon as the parameter $n$ of the representation $(V_n, \chi_n)$ (notation as in the proof of Proposition 5.1) is even, (39) implies that

$$\chi_n \left( \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} \right)$$

acts trivially on all elements of $V_n$. Hence $\chi_n$ descends to a representation of $\text{PSL}_2(\mathbb{R})$.

We remark that, by the same reasoning, the representations $(V_n, \chi_n)$ with odd parameter $n$ do not descend to $\text{PSL}_2(\mathbb{R})$.

For the next class of examples of representations with non-expanding cusp monodromies we recall the signature of Fuchsian groups and its associated presentations. For details we refer to [2, §10.4].

Let $\Gamma$ be a non-elementary geometrically finite Fuchsian group, and let $g$ be the genus of $\Gamma \setminus \mathbb{H}$. Further let $r$ be the number of conjugacy classes of maximal elliptic cyclic subgroups of $\Gamma$ (or, equivalently, the number of conical points of $\Gamma \setminus \mathbb{H}$), and suppose that the orders of these subgroups are $m_1, \ldots, m_r$, respectively. Let $s$ be the number of conjugacy classes of maximal parabolic cyclic subgroups in $\Gamma$ (or, equivalently, the number of cusps of $\Gamma \setminus \mathbb{H}$), and let $t$ be the number of conjugacy classes of maximal boundary hyperbolic cyclic subgroups (or, equivalently, the number of funnels of $\Gamma \setminus \mathbb{H}$). The **signature** of $\Gamma$ is then the tuple

$$(g; m_1, \ldots, m_r; s; t). \quad (41)$$

Each parameter is a non-negative integer, and $m_j \geq 2$ for all $j \in \{1, \ldots, r\}$. If $\Gamma$ does not have elliptic elements, then we write $(g; 0; s; t)$.

For any $g, r, s, t \in \mathbb{N}$ and any $m_j \in \mathbb{N}_{\geq 2}$ for $j \in \{1, \ldots, r\}$, there exists a non-elementary geometrically finite Fuchsian group $\Gamma$ with signature $(g; m_1, \ldots, m_r; s; t)$.
if and only if

\[ 2g - 2 + s + t + \sum_{j=1}^{r} \left(1 - \frac{1}{m_j}\right) > 0. \]

See [2, Theorem 10.4.2].

If \( \Gamma \) is a geometrically finite Fuchsian group with signature \((g; m_1, \ldots, m_r; s; t)\), then \( \Gamma \) has a presentation of the form

\[ \Gamma = \left\langle a_1, b_1, \ldots, a_g, b_g, e_1, \ldots, e_r, p_1, \ldots, p_s, h_1, \ldots, h_t \mid e_1^{m_1} = \cdots = e_r^{m_r} \right. \]

\[ \left. = \left( \prod_{j=1}^{g} [a_j, b_j] \right) e_1 \cdots e_r p_1 \cdots p_s h_1 \cdots h_t = 1 \right\rangle, \]

where the elements \( a_1, b_1, \ldots, a_g, b_g \) are hyperbolic (being responsible for the genus of \( \Gamma \)), \( e_1, \ldots, e_r \) are elliptic (generators of the maximal elliptic cyclic subgroups), \( p_1, \ldots, p_s \) are parabolic (generators of the maximal parabolic subgroups), \( h_1, \ldots, h_t \) are hyperbolic elements (generators of the maximal boundary hyperbolic cyclic subgroups).

**Proposition 5.3** Let \( \Gamma \) be a geometrically finite Fuchsian group with at least one parabolic element. Then there exists a finite-dimensional representation of \( \Gamma \) with non-expanding cusp monodromy that is neither unitary at cusps nor the restriction of a representation of \( SL_2(\mathbb{R}) \) nor decomposes into a direct sum of representations of smaller dimensions.

Throughout the proof we will use ‘direct sum’ to mean a ‘direct sum of representations of smaller dimensions.’

**Proof** Let \( \Gamma \) be a geometrically finite non-uniform Fuchsian group with at least one parabolic element, and suppose that the signature of \( \Gamma \) is \((g; m_1, \ldots, m_r; s; t)\). The presence of a parabolic element in \( \Gamma \) shows that \( s \geq 1 \). Further, the signature of \( \Gamma \) has to satisfy at least one of the following properties:

(a) \( g \geq 1 \),
(b) \( r \geq 1 \) and \( m_j \neq 3 \) for some \( j \in \{1, \ldots, r\} \),
(c) \( r \geq 1 \) and \( m_j = 3 \) for some \( j \in \{1, \ldots, r\} \),
(d) \( s \geq 2 \),
(e) \( t \geq 1 \), or
(f) the signature of \( \Gamma \) is \((0; 0; 1; 0)\).

To prove the proposition we will construct an (explicit) example of a representation

\[ \varrho : \Gamma \to GL(V) \]
on some finite-dimensional vector space $V$ such that $\varrho$ has non-expanding cusp monodromy but is neither unitary at cusps nor the restriction of a representation of $\text{SL}_2(\mathbb{R})$ nor decomposes into a direct sum. The construction of $\varrho$ will depend on whether $\Gamma$ satisfies (a), (b), (c), (d), (e) or (f). (In case that $\Gamma$ satisfies more than one of these properties, our construction yields more than one example.)

We first discuss the case that $\Gamma$ satisfies (f). Then $\Gamma$ is freely generated by a single parabolic element, say $\Gamma = \langle p \rangle$. The representation
\[ \varrho : \Gamma \rightarrow \text{GL}_2(\mathbb{R}), \]
determined by
\[ \varrho(p) := \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, \]
has non-expanding cusp monodromy but is not unitary (at the cusp). Since irreducible two-dimensional representations of $\text{SL}_2(\mathbb{R})$ do not restrict to $\text{PSL}_2(\mathbb{R})$ and hence a fortiori not to $\Gamma$ (see Remark 5.2), the representation $\varrho$ is not a restriction of an irreducible representation of $\text{SL}_2(\mathbb{R})$. Further, $\varrho$ is not a direct sum of two one-dimensional representations because $\varrho(p)$ does not diagonalize. A fortiori, $\varrho$ is not the restriction of a representation of $\text{SL}_2(\mathbb{R})$. Hence, $\varrho$ is a representation of $\Gamma$ of the form claimed.

We now suppose that $\Gamma$ satisfies any of the properties (a)–(e). Then let
\[ \Gamma = \left\langle a_1, b_1, \ldots, a_g, b_g, e_1, \ldots, e_r, p_1, \ldots, p_s, h_1, \ldots, h_t \mid e_1^{m_1} = \cdots = e_r^{m_r} \right\rangle \]
be a presentation of $\Gamma$ that realizes its signature. Thus, the elements $a_1, b_1, \ldots, a_g, b_g$ are hyperbolic, $e_1, \ldots, e_r$ are elliptic, $p_1, \ldots, p_s$ are parabolic, $h_1, \ldots, h_t$ are boundary hyperbolic elements. Throughout we assume without loss of generality that
\[ p_1 = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}. \]
Obviously, it suffices to define $\varrho$ on the set of generators
\[ \mathcal{G} := \{ a_1, b_1, \ldots, a_g, b_g, e_1, \ldots, e_r, p_1, \ldots, p_s, h_1, \ldots, h_t \} \]
for $\Gamma$, checking that it obeys all relations in the presentation of $\Gamma$ from above, and then extend it (in the unique possible way) to all of $\Gamma$ enforcing homomorphy.

We suppose first that $\Gamma$ satisfies (a). In this case, we set $V := \mathbb{C}^2$. We fix $\lambda := e^{i \pi / 2}$, $\mu := e^{i \pi / 6}$. We set
\[ \varrho(a_1) := \begin{pmatrix} \lambda & 1 \\ 0 & \lambda^{-1} \end{pmatrix}, \quad \varrho(b_1) := \begin{pmatrix} \mu & 1 \\ 0 & \mu^{-1} \end{pmatrix}, \]
\[ \varrho(p_1) := \begin{pmatrix} 1 & \lambda \mu (\lambda^{-1} - \lambda + \mu - \mu^{-1}) \\ 0 & 1 \end{pmatrix}. \]
and
\[ \varrho(g) := \text{id}_{C^2} \quad \text{for all } g \in G \setminus \{a_1, b_1, p_1\}. \]

Then the relations in the presentation of $\Gamma$ are satisfied. Extending $\varrho$ to all of $\Gamma$ by enforcing homomorphy then yields a well-defined finite-dimensional representation of $\Gamma$. From (42) we read off that $\varrho$ has non-expanding cusp monodromy. Further, since
\[ \lambda \mu (\lambda^{-1} - \lambda + \mu - \mu^{-1}) = \frac{1}{2} (i + \sqrt{3}) \neq 0, \]
the representation $\varrho$ is not unitary at cusps. Since $\varrho$ is two-dimensional, it cannot be the restriction of an irreducible representation of $\text{SL}_2(\mathbb{R})$ (see Remark 5.2). Further, since $\varrho(p_1)$ is not diagonalizable, $\varrho$ is not the direct sum of two one-dimensional representations. Therefore $\varrho$ is of the type of claimed.

We now suppose that $\Gamma$ satisfies (b), (d) or (e). In each case we let $(V_2, \chi_2)$ be the representation of $\text{SL}_2(\mathbb{R})$ from (39) with $n = 2$ and let $\chi'_2$ be the restriction of $\chi_2$ to $\Gamma$. We set $V := V_2$ and will define $\varrho$ by modifying $\chi'_2$ as explained in what follows.

(i) If $\Gamma$ satisfies (b) then let $\tau$ be an $m_1$-th root of unity such that $\tau^3 \neq 1$. We set
\[ \varrho(e_1) := \tau^{-1} \cdot \chi'_2(e_1), \quad \varrho(p_1) := \tau \cdot \chi'_2(p_1), \]
and
\[ \varrho(g) := \chi'_2(g) \quad \text{for } g \in G \setminus \{e_1, p_1\}. \]

(ii) If $\Gamma$ satisfies (d) then we pick $\tau \in \mathbb{C}$ such that $|\tau| = 1$ and $\tau^3 \neq 1$. We set
\[ \varrho(p_1) := \tau \cdot \chi'_2(p_1), \quad \varrho(p_2) := \tau^{-1} \cdot \chi'_2(p_2), \]
and
\[ \varrho(g) := \chi'_2(g) \quad \text{for } g \in G \setminus \{p_1, p_2\}. \]

(iii) If $\Gamma$ satisfies (e) then we again pick $\tau \in \mathbb{C}$ such that $|\tau| = 1$ and $\tau^3 \neq 1$. We set
\[ \varrho(p_1) := \tau \cdot \chi'_2(p_1), \quad \varrho(h_1) := \tau^{-1} \cdot \chi'_2(h_1), \]
and
\[ \varrho(g) := \chi'_2(g) \quad \text{for } g \in G \setminus \{p_1, h_1\}. \]

In each case, one easily checks that the relations in the presentation of $\Gamma$ are satisfied. Further, (40) and $\varrho(p_1) = \tau \chi'_2(p_1)$ imply that $\varrho(p_1)$ is not unitary. Hence $\varrho$ is not
unitary at cusps. Since $\varrho(p_1)$ consists of a single Jordan block (see (40)), $\varrho$ does not decompose into a direct sum.

Finally, if $\varrho$ was the restriction of an irreducible representation of $\text{SL}_2(\mathbb{R})$, then $\varrho$ would be equivalent to $\chi'_2$. Then there would exist an intertwiner $\phi : V_2 \to V_2$ such that $\phi \circ \varrho = \chi'_2 \circ \phi$. This would imply

$$1 = \det(\chi'_2(p_1)) = \det(\varrho(p_1)) = \tau^3 \det(\chi'_2(p_1)) = \tau^3,$$

which contradicts $\tau^3 \neq 1$.

If $\Gamma$ satisfies (c) then we perform a similar construction but start with the representation $(V_4, \chi_4)$ of $\text{SL}_2(\mathbb{R})$ (from (39)). Let $\chi'_4$ denote its restriction to $\Gamma$, and let $\tau$ be a third root of unity such that $\tau^5 \neq 1$. We set

$$\varrho(e_1) := \tau^{-1} \cdot \chi'_4(e_1), \quad \varrho(p_1) := \tau \cdot \chi'_4(p_1),$$

and

$$\varrho(g) := \chi'_4(g) \quad \text{for } g \in G \setminus \{e_1, p_1\}.$$

It follows that $\varrho$ is not unitary at cusps. Since $\varrho(p_1)$ again consists of a single Jordan block, it does not decompose into a direct sum. Further, if $\varrho$ was a restriction of an irreducible representation of $\text{SL}_2(\mathbb{R})$, then it would be equivalent to $\chi'_4$ and hence, analogously, to above

$$1 = \det(\chi'_4(p_1)) = \det(\varrho(p_1)) = \tau^5 \det(\chi'_4(p_1)) = \tau^5,$$

which is a contradiction. This completes the proof. \(\Box\)

We end this section with the proof that for representations the property of non-expanding cusp monodromy is stable under induction.

**Proposition 5.4** Let $\Gamma$ be a Fuchsian group, $\Lambda$ a subgroup of $\Gamma$ of finite index and $\eta : \Lambda \to \text{GL}(V)$ a finite-dimensional representation with non-expanding cusp monodromy. Then the induced representation $\chi := \text{Ind}_\Lambda^\Gamma \eta$ of $\Gamma$ has non-expanding cusp monodromy.

**Proof** Let $\tilde{\eta} : \Gamma \to \text{End}(V)$ be defined by

$$\tilde{\eta}(g) := \begin{cases} \eta(g) & \text{if } g \in \Lambda \\ 0 & \text{otherwise}. \end{cases}$$

Let $n := [\Gamma : \Lambda]$, and let $R = \{h_1, \ldots, h_n\}$ be a complete set of representatives in $\Gamma$ for the right cosets $\Lambda \setminus \Gamma$ (thus, $\Gamma = \bigcup_{j=1}^n \Lambda h_j$ as a disjoint union). Then $\chi$ is (equivalent to) the representation on $V^n \cong \bigoplus_{j=1}^n V$ given by

$$\chi(g) = \left(\tilde{\eta}(h_i g h_j^{-1})\right)_{i,j=1,\ldots,n} \quad (g \in \Gamma).$$

(43)
Let \( p \in \Gamma \) be parabolic. Since \( R \) is a complete set of representatives for \( \Lambda \setminus \Gamma \), for each \( i_0 \in \{1, \ldots, n\} \) there exists exactly one \( j_0 \in \{1, \ldots, n\} \), and for each \( j_0 \in \{1, \ldots, n\} \) there exists exactly one \( i_0 \in \{1, \ldots, n\} \) such that \( h_{i_0} p h_{j_0}^{-1} \in \Lambda \). In other words, in the presentation (43), the endomorphism \( \chi(p) \) is given by a ‘permutation matrix’ where instead of the entry 1 (at position \((i, j)\)) we have \( \eta(h_i p h_j^{-1}) \). It follows that there exists \( k \in \mathbb{N} \) (indeed \( k \leq n! \)) such that

\[
\chi(p^k) = \chi(p)^k = \text{diag} \left( \eta(h_1 p^k h_1^{-1}), \ldots, \eta(h_n p^k h_n^{-1}) \right).
\]

Since all eigenvalues of the endomorphisms \( \eta(h_j p^k h_j^{-1}) \), \( j = 1, \ldots, n \), are of absolute value 1, so are all eigenvalues of \( \chi(p)^k \), and hence all eigenvalues of \( \chi(p) \). Thus, \( \chi \) has non-expanding cusp monodromy.

### 6 No convergence beyond non-expanding cusp monodromy

In the preceding sections we showed that if \( \chi \) is a representation with non-expanding cusp monodromy, then the infinite product (8) used for the definition of the \( \chi \)-twisted Selberg zeta function converges absolutely in a half space, and has a meromorphic continuation to all of \( \mathbb{C} \). In this section we will discuss the natural question if these results hold as well for some representations \( \chi \) not having non-expanding cusp monodromy.

The infinite product (8) contains the product over the primitive periodic geodesics or, equivalently, the \( \Gamma \)-conjugacy classes of the primitive hyperbolic elements in \( \Gamma \), which does not have a natural order. Therefore, absolute convergence is arguably the most natural notion of convergence when considering (8).

We will show that as soon as \( \chi \) does not have non-expanding cusp monodromy and the infinite product (8) is not void then the latter does not converge absolutely, and hence the notion of twisted Selberg zeta function has no meaning (within the realms of absolute convergence). This observation shows that the property of non-expanding cusp monodromy is indeed the weakest requirement on twisting representations we can allow for a theory of Selberg zeta functions.

**Proposition 6.1** Let \( \Gamma \) be a geometrically finite Fuchsian group with at least one hyperbolic element, and let \( \chi : \Gamma \to \text{GL}(V) \) be a finite-dimensional representation of \( \Gamma \) that does not have non-expanding cusp monodromy. Then the infinite product (8) does not converge absolutely for any \( s \in \mathbb{C} \).

**Proof** Since \( \chi \) does not have non-expanding cusp monodromy, there exists a parabolic element \( p \in \Gamma \) such that not all eigenvalues of \( \chi(p) \) are of absolute value 1. Without loss of generality (by possibly conjugating \( \Gamma \), and taking the inverse of \( p \)) we may assume that

\[
p = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}
\]
and that $\chi(p)$ has an eigenvalue $\lambda$ with $|\lambda| > 1$. Further, for readability, we assume that $\chi(p)$ has a single Jordan block. The general case of several Jordan blocks is easily handled by a straightforward extension of the argumentation that follows.

By hypothesis, $\Gamma$ has at least one hyperbolic element, hence the associated hyperbolic surface $X = \Gamma \backslash \mathbb{H}$ has at least one (primitive) periodic geodesic, and the infinite product (8) is non-empty. Let

$$ h = \begin{bmatrix} a & b \\ c & d \end{bmatrix} $$

be a hyperbolic element in $\Gamma$. In what follows we will show that $(hp^n)_{n \in \mathbb{N}}$ contains a subsequence of hyperbolic elements such that even if the infinite product (8) is restricted to these hyperbolic elements (or rather its underlying primitive elements), it does not converge. Geometrically, it means that we restrict the considerations to a periodic geodesic winding arbitrarily often around the cusp associated to $p$. The weights that $\chi$ assigns to these geodesics are so large that convergence of (8) is prevented.

By the hypothesis above on $p$, the point $\infty$ is cuspidal and hence the element $c$ in $h$ does not vanish. A straightforward calculation shows that for all $n \in \mathbb{N}$ we have

$$ \text{tr}(hp^n) = \text{tr} h + cn. $$

This shows that there exists $n_0 \in \mathbb{N}$ such that the elements of $\{hp^n \mid n \geq n_0\}$ are all hyperbolic and pairwise non-conjugate in $\Gamma$.

As discussed in Sect. 3, if the infinite product (8) converges absolutely for $s \in \mathbb{C}$ then

$$ \sum_{[g] \in [\Gamma]_h} \frac{1}{m(g)} \frac{N(g)^{-s}}{1 - N(g)^{-1}} \text{tr} \chi(g) $$

converges absolutely. In particular, for $n_0 \in \mathbb{N}$ sufficiently large, the sum

$$ \Sigma(s) := \sum_{n \geq n_0} \frac{1}{m(hp^n)} \frac{N(hp^n)^{-\text{Re}s}}{1 - N(hp^n)^{-1}} |\text{tr} \chi(hp^n)| $$

(44)

converges for these $s \in \mathbb{C}$. In what follows we show that the sum (44) does not converge.

To that end, we start by estimating $N(hp^n)$ and $m(hp^n)$ for $n \in \mathbb{N}$. The eigenvalues of $hp^n$ are

$$ \mu_{\pm} = \frac{a + d + cn \pm \sqrt{(a + d + cn)^2 - 4}}{2}. $$

Thus (see (6))

$$ N(hp^n) \asymp n^2 $$

(45)
with implied constants independent of $n$. Let
\[ \varepsilon_0 := \min\{N(h_0) \mid h_0 \in \Gamma \text{ hyperbolic}\}. \]

We note that this minimum indeed exists. Suppose that $h_n$ is the primitive hyperbolic element underlying $ph^n$, and let $m := m(hp^n)$. Hence $h^m_n = hp^n$ and therefore
\[ \varepsilon^m_0 \leq N(h_n)^m = N(h^m_n) = N(hp^n) \asymp n^2. \]

It follows that
\[ m(hp^n) = m \ll \frac{1}{\log n} \tag{46} \]
with implied constant independent of $n$. Now let $d := \dim V$ and note that there exist $a_0, \ldots, a_{d-1} \in \mathbb{C}$ such that for all $n \in \mathbb{N}$ sufficiently large we have
\[ \text{tr} \chi(hp^n) = \lambda^n \left(a_0 + a_1 n + \cdots + a_{d-1} n^{d-1}\right) \gg n^{d-1} |\lambda|^n \tag{47} \]
with implied constant independent of $n$. Using (45)–(47) in (44) we find that
\[ \Sigma(s) \gg \sum_{n \geq n_0} \frac{1}{\log n} \frac{n^{-2 \Re s}}{1 - n^{-2} n^{d-1} |\lambda|^n} \]
\[ \gg \sum_{n \geq n_0} n^k |\lambda|^n \]
for some $k \in \mathbb{R}$. Since $|\lambda| > 1$, the last sum diverges. This completes the proof. \hfill \Box

7 Venkov–Zograf factorization formulas

Venkov and Zograf [59] showed that Selberg zeta functions for geometrically finite Fuchsian groups twisted with unitary finite-dimensional representations satisfy certain factorization formulas. The purpose of this section is to provide, with Theorem 7.1 below, these factorization formulas for twists with non-expanding cusp monodromy.

To be precise, the proof of the factorization formulas in [59] takes advantage of the Selberg trace formula, which restricts the proof method from [59] to cofinite Fuchsian groups and unitary representations. However, the simplification (which essentially consists in omitting some unnecessary steps) sketched in [58, Theorem 7.2] applies without changes to non-cofinite Fuchsian groups as well, even though [58, Theorem 7.2] is stated for cofinite Fuchsian groups only.

In combination with Proposition 5.4 the latter proof applies also to all finite-dimensional representations with non-expanding cusp monodromy, again without any changes. For the convenience of the reader we provide a complete proof.

**Theorem 7.1** Let $\Gamma$ be a geometrically finite Fuchsian group, and $\Lambda$ a subgroup of $\Gamma$ of finite index. Let $\chi : \Gamma \rightarrow \text{GL}(V)$ be a finite-dimensional representation with non-expanding cusp monodromy.
(i) If \( \chi = \bigoplus_{j=1}^m \chi_j \) decomposes into a direct (finite) sum of representations \( \chi_j \) of \( \Gamma \) with non-expanding cusp monodromy, then

\[
Z(s, \Gamma, \chi) = \prod_{j=1}^m Z(s, \Gamma, \chi_j).
\]

(ii) If \( \eta: \Lambda \to \text{GL}(V) \) is a representation with non-expanding cusp monodromy, then

\[
Z(s, \Lambda, \eta) = Z(s, \Gamma, \text{Ind}\_\Lambda^\Gamma \eta).
\]

Recall from Proposition 5.4 that \( \text{Ind}\_\Lambda^\Gamma \eta \) has indeed non-expanding cusp monodromy.

(ii) If \( \Lambda \) is normal in \( \Gamma \), then

\[
Z(s, \Lambda, 1) = \prod_{\eta \in \hat{\Lambda \backslash \Gamma}} Z(s, \Gamma, \eta)^{\dim \eta},
\]

where \( \hat{\Lambda \backslash \Gamma} \) denotes the unitary dual of \( \Lambda \backslash \Gamma \).

**Proof** The proof of (i) follows immediately from (8) in the domain of convergence of the infinite products and extends to all of \( \mathbb{C} \) by uniqueness of meromorphic continuations. In the case that \( \Gamma \) admits a strict transfer operator approach, an alternative transfer operator based proof of (i) is possible and straightforward since the transfer operators decompose according to the representation. We refer to [50, Section 6] for details.

Further, (ii) follows immediately from (i) and (ii) by considering the decomposition of the induced representation \( \text{Ind}\_\Lambda^\Gamma \eta \) into irreducibles.

It remains to prove (ii). To that end we show that the logarithms of the two Selberg zeta functions coincide on the domain on which both of them are given by the infinite products in (8). Uniqueness of meromorphic continuations then completes the proof. We start with a few preparations.

We define, as in the proof of Proposition 5.4, the map \( \tilde{\eta}: \Gamma \to \text{End}(V) \) by

\[
\tilde{\eta}(g) := \begin{cases} \eta(g) & \text{if } g \in \Lambda \\ 0 & \text{otherwise.} \end{cases}
\]

For \( g, h \in \Gamma \) we write \( g \sim_\Gamma h \) or \( g \sim_\Lambda h \) if \( g \) and \( h \) are \( \Gamma \)-conjugate or \( \Lambda \)-conjugate, respectively.

Let \( \chi := \text{Ind}\_\Lambda^\Gamma \eta \), and let \( R \) be a complete set of representatives for the right cosets \( \Lambda \backslash \Gamma \). Thus,

\[
\Gamma = \bigcup_{h \in R} \Lambda h \quad (\text{disjoint union}).
\]
Note that $R$ is finite. For any $g \in \Gamma$ we have
\[
\text{tr } \chi(g) = \sum_{h \in R} \text{tr } \tilde{\eta}(hgh^{-1})
\]  
(48)
by Frobenius formula (also known as Mackey formula).

Let $p_0 \in \Gamma$ be $\Gamma$-primitive hyperbolic, and consider
\[
\Lambda(p_0) := \{gp_0^kg^{-1} | g \in \Gamma, k \in \mathbb{N}\} \cap \Lambda.
\]

Then [59, Lemma 2.3] shows that whenever $gp_0^kg^{-1} \in \Lambda(p_0)$ for some $g \in \Gamma$, $k \in \mathbb{N}$ then there exists a $\Lambda$-primitive element $q_0 \in \Lambda$ and a unique $e_0 \in \mathbb{N}$ such that $q_0 \sim \Gamma p_0^{e_0}$ and $e_0$ divides $k$. Thus, $q_0 \in \Lambda(p_0)$.

By slight abuse of notation, let $[\Lambda(p_0)]_p$ denote the set of $\Lambda$-conjugacy classes of $\Lambda$-primitive hyperbolic elements in $\Lambda(p_0)$. By [59, Lemmas 2.1-2.2], $[\Lambda(p_0)]_p$ is finite. Let
\[
R(p_0) := \{q_1, \ldots, q_M\}
\]
be a complete set of representatives of $[\Lambda(p_0)]_p$ in $\Lambda$ (thus, $q_1, \ldots, q_M$ are $\Lambda$-primitive hyperbolic and pairwise non-conjugate in $\Lambda$, and each element of $[\Lambda(p_0)]_p$ has a representative in $R(p_0)$). For $j \in \{1, \ldots, M\}$ let $e_j \in \mathbb{N}$ such that $q_j \sim \Gamma p_0^{e_j}$, and let
\[
R_j := \{h \in R | q_j \sim \Lambda h p_0^{e_j} h^{-1}\}.
\]
By [59, Lemma 2.1], $|R_j| = e_j$. It follows that for all $k \in \mathbb{N}$ with $e_j \mid k$ we have
\[
e_j \text{tr } \eta(q_j^{k/e_j}) = \sum_{h \in R_j} \text{tr } \eta(hp_0^k h^{-1}).
\]

Thus, for all $k \in \mathbb{N}$ it follows
\[
\sum_{j=1}^{M} \frac{e_j \text{tr } \eta(q_j^{k/e_j})}{e_j \mid k} = \sum_{h \in R} \text{tr } \tilde{\eta}(hp_0^k h^{-1}) = \text{tr } \chi(p_0^k).
\]

Hence
\[
\sum_{k=1}^{\infty} \frac{1}{k} \frac{N(p_0^k)^{-s}}{1 - N(p_0^k)} \text{tr } \chi(p_0^k) = \sum_{k=1}^{\infty} \sum_{j=1}^{M} \frac{e_j}{k} \frac{N(p_0^k)^{-s}}{1 - N(p_0^k)} \text{tr } \eta(q_j^{k/e_j})
\]
\[
= \sum_{j=1}^{M} \sum_{m=1}^{\infty} \frac{e_j}{e_j m} \frac{N(p_0^{e_j m})^{-s}}{1 - N(p_0^{e_j m})} \text{tr } \eta(q_j^m)
\]
\[
= \sum_{j=1}^{M} \sum_{m=1}^{\infty} \frac{1}{m} \frac{N(q_j^m)^{-s}}{1 - N(q_j^m)} \text{tr } \eta(q_j^m).
\]
Thus, for \( \Re s \) sufficiently large, 

\[
\log Z(s, \Gamma, \chi) = - \sum_{[p_0] \in \Gamma_p} \sum_{k=1}^{\infty} \frac{1}{k} \frac{N(p_k^0)^{-s}}{1 - N(p_k^0)^{-1}} \text{tr} \chi(p_k^0) 
\]

\[
= - \sum_{[p_0] \in \Gamma_p} \sum_{q \in \Lambda(p_0)} \sum_{m=1}^{\infty} \frac{1}{m} \frac{N(q^m)^{-s}}{1 - N(q^m)^{-1}} \text{tr} \eta(q^m) 
\]

\[
= - \sum_{[g] \in \Lambda_p} \sum_{m=1}^{\infty} \frac{1}{m} \frac{N(g^m)^{-s}}{1 - N(g^m)^{-1}} \text{tr} \eta(g^m) 
\]

\[
= \log Z(s, \Lambda, \eta). 
\]

This completes the proof. \( \square \)

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Appendix A: Meromorphic continuation of the derivatives of the Lerch transcendent

In Sect. 4.7 we used the Lerch transcendent

\[
\Phi(s, \lambda, w) = \sum_{n=0}^{\infty} \frac{\lambda^n}{(n + w)^s} 
\]

and its partial derivatives

\[
\Phi_m(s, \lambda, w) = \frac{1}{m!} \frac{\partial^m \Phi}{\partial \lambda^m}(s, \lambda, w) = \sum_{n=m}^{\infty} \binom{n}{m} \frac{\lambda^{n-m}}{(n + w)^s} \quad (m \in \mathbb{N}_0) 
\]

for the proof of the meromorphic continuability of the transfer operator families. In particular we took advantage of the convergence of \( \Phi_m, m \in \mathbb{N}_0 \), their extensions to meromorphic or even entire functions to all of \( \mathbb{C} \) in the variable \( s \) for any \( w \in \mathbb{C} \setminus (-\mathbb{N}_0) \), and their continuity as a function of the two variables \( (s, w) \) for any
$\lambda \in \mathbb{C}, |\lambda| = 1$. Proposition A.1 below contains a precise statement of the properties we used.

For $\Phi = \Phi_0$, proofs of these properties can be found in the literature. For $\Phi_m$, $m > 0$, they can be deduced from those of $\Phi$ (at least in the case $\lambda \neq 1$) or shown in analogy to those of $\Phi$. For the convenience of the reader, we provide a complete proof of Proposition A.1.

For this proof we take an approach via Mellin transforms and functional equations. A descent reference for this approach is Zagier’s paper [60].

**Proposition A.1** Let $\lambda \in \mathbb{C}, |\lambda| = 1, m \in \mathbb{N}_0$. Then the infinite sum $\Phi_m(\cdot, \lambda, \cdot)$ in (50) is convergent on

$$\{s \in \mathbb{C} \mid \Re s > m + 1\} \times (\mathbb{C} \setminus (-\mathbb{N}_0)),$$

defines there a continuous function, and for any fixed $w \in \mathbb{C} \setminus (-\mathbb{N}_0)$, the map $\Phi_m(\cdot, \lambda, w)$ is holomorphic on

$$\{s \in \mathbb{C} \mid \Re s > m + 1\}.$$

Furthermore, $\Phi_m(\cdot, \lambda, \cdot)$ is holomorphic on

$$\{s \in \mathbb{C} \mid \Re s > m + 1\} \times (\mathbb{C} \setminus (-\mathbb{N}_0)).$$

For $\lambda \neq 1$, the map $\Phi_m(\cdot, \lambda, \cdot)$ extends to a holomorphic function on

$$\mathbb{C} \times (\mathbb{C} \setminus (-\mathbb{N}_0)).$$

For $\lambda = 1$, it extends to a meromorphic function on $\mathbb{C} \times (\mathbb{C} \setminus (-\mathbb{N}_0))$. For any fixed $w \in \mathbb{C} \setminus (-\mathbb{N}_0)$, the poles of $\Phi_m(\cdot, 1, w)$ are contained in $\{1, \ldots, m + 1\}$, all poles are simple, and $m + 1$ is a pole.

If a continuous extension of the complex logarithm to $\mathbb{C} \setminus (-\mathbb{N}_0)$ is fixed, then the meromorphic extendability in $s$ generalizes to all fixed $w \in \mathbb{C} \setminus (-\mathbb{N}_0)$, and the meromorphic extension is continuous as a function of the two variables $(s, w)$.

**Proof** Throughout let $\lambda \in \mathbb{C}, |\lambda| = 1$, and let $m \in \mathbb{N}_0$. Definitions (49) and (50) clearly show that $\Phi_m(\cdot, \lambda, \cdot)$ converges on

$$\mathcal{D}_m := \{s \in \mathbb{C} \mid \Re s > m + 1\} \times (\mathbb{C} \setminus (-\mathbb{N}_0)),$$

defines a continuous function on this domain, and for any fixed $w \in \mathbb{C} \setminus (-\mathbb{N}_0)$, the map $\Phi_m(\cdot, \lambda, w)$ is holomorphic on $\{s \in \mathbb{C} \mid \Re s > m + 1\}$. Furthermore, the map $\Phi_m(\cdot, \lambda, \cdot)$ is holomorphic on

$$\{s \in \mathbb{C} \mid \Re s > m + 1\} \times (\mathbb{C} \setminus (-\infty, 0]).$$
For the proof of the claimed meromorphic continuability we first note that for all \( k \in \mathbb{N}_0 \) and all \((s, w) \in D_0\) we have

\[
\Phi(s, \lambda, w) = \lambda^k \Phi(s, \lambda, w + k) + \sum_{n=0}^{k-1} \frac{\lambda^n}{(n+w)^s}.
\]

Moreover, for every \( n \in \mathbb{N}_0 \), the map

\[
\mathbb{C} \times (\mathbb{C} \setminus (-\infty, 0]) \to \mathbb{C}, \quad (s, w) \mapsto \frac{\lambda^n}{(n+w)^s}
\]

is holomorphic.

Combining this functional equation with the definition (50) of \( \Phi_m \) as a partial derivative of \( \Phi \) yields that for all \( k \in \mathbb{N}_0 \) and all \((s, w) \in D_m\),

\[
\Phi_m(s, \lambda, w) = \sum_{n=0}^{m} \binom{k}{n} \lambda^{k-n} \Phi_{m-n}(s, \lambda, w + k) + \sum_{n=0}^{k-1} \binom{n}{m} \frac{\lambda^{n-m}}{(n+w)^s}. \tag{51}
\]

Suppose for a moment that we have already shown (which will be done further below) that for any \( v \in \mathbb{C} \) with \( \text{Re} \, v > 0 \) and any \( p \in \mathbb{N}_0 \), the map \( \Phi_p(\cdot, \lambda, v) \) extends, for \( \lambda \neq 1 \), to an entire function, and, for \( \lambda = 1 \), to a meromorphic function on all of \( \mathbb{C} \) all of whose poles are contained in \( \{1, \ldots, p+1\} \), are (at most) simple, and \( p+1 \) is indeed a pole. Suppose further that we have already shown that these extensions are meromorphic in the variables \((s, v)\) on

\[
(\mathbb{C} \setminus \{\text{poles}\}) \times \{v \in \mathbb{C} \mid \text{Re} \, v > 0\}.
\]

Applying (51) with sufficiently large (and varying) \( k \in \mathbb{N}_0 \) proves the statement of the proposition for \( w \) restricted to \( \mathbb{C} \setminus (-\infty, 0] \). Fixing a continuous extension of the complex logarithm to \( \mathbb{C} \setminus (-\mathbb{N}_0) \) then completes the full proof of the proposition.

Thus it remains to establish the proposition for the domain of the variable \( w \) restricted to \( \text{Re} \, w > 0 \). We start by expressing \( \Phi_m \) by a Mellin transform of a certain highly regular function with good decay properties. Throughout let \( w \in \mathbb{C} \) with \( \text{Re} \, w > 0 \), and \( s \in \mathbb{C} \) with \( \text{Re} \, s > m + 1 \). The integral representation

\[
\Gamma(r) = \int_0^\infty u^{r-1} e^{-u} du, \quad \text{Re} \, r > 0
\]

of the Gamma function yields for all \( n \in \mathbb{N}_0 \) that

\[
\frac{1}{(n+w)^r} = \frac{1}{\Gamma(r)} \int_0^\infty e^{-t(n+w)} t^{r-1} dt
\]
(initially only for $w \in \mathbb{R}_{>0}$ and then by holomorphy and identity theorem also for $w \in \mathbb{C}, \text{Re } w > 0$). It follows that (note that $|\lambda e^{-t}| < 1$ for $t > 0$)

$$\Phi_m(s, \lambda, w) = \sum_{n=m}^{\infty} \binom{n}{m} \frac{\lambda^{n-m}}{(n+w)^s}$$

$$= \frac{1}{\Gamma(s)} \sum_{n=m}^{\infty} \binom{n}{m} \lambda^{n-m} \int_0^\infty t^{s-1} e^{-t(n+w)} dt$$

$$= \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} e^{-(m+w)t} \sum_{n=m}^{\infty} \binom{n}{m} (\lambda e^{-t})^{n-m} dt$$

$$= \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} \frac{e^{-(w+1)t}}{(1-\lambda e^{-t})^{m+1}} dt$$

$$= \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} \frac{e^{-(w-1)t}}{(e^{-t} - \lambda)^{m+1}} dt.$$

Let $\varphi_{\lambda, w} : \mathbb{R} \to \mathbb{C}$,

$$\varphi_{\lambda, w}(t) := \frac{e^{-(w-1)t}}{(e^{-t} - \lambda)^{m+1}},$$

and let $\tilde{\varphi}_{\lambda, w} : \{ s \in \mathbb{C} \mid \text{Re } s > m + 1 \} \to \mathbb{C}$,

$$\tilde{\varphi}_{\lambda, w}(s) := \int_0^\infty t^{s-1} \varphi_{\lambda, w}(t) dt$$

denote the Mellin transform of $\varphi_{\lambda, w}$. Thus

$$\Phi_m(s, \lambda, w) = \Gamma(s)^{-1} \tilde{\varphi}_{\lambda, w}(s).$$

Since $\text{Re } w > 0$, one easily sees that $\varphi_{\lambda, w}$ is of rapid decay as $t \to \infty$. Moreover, as $t \to 0$, the function $\varphi_{\lambda, w}$ has an asymptotic expansion

$$\varphi_{\lambda, w}(t) \sim \sum_{n=-(m+1)}^{\infty} a_n(\lambda, w) t^n$$

with suitable coefficients $a_n(\lambda, w) \in \mathbb{C}$. For $\lambda \neq 1$, $\varphi_{\lambda, w}$ is smooth at 0, and therefore

$$a_n(\lambda, w) = 0 \quad \text{if } \lambda \neq 1 \text{ and } n \in \{-(m+1), \ldots, -1\}.$$

For $\lambda = 1$, $\varphi_{\lambda, w}$ has a pole of order $m + 1$ at 0.
For any $N \in \mathbb{Z}$, $N > -(m + 1)$, we find

$$
\tilde{\phi}_{\lambda, w}(s) = \int_0^\infty t^{s-1} \left( \varphi_{\lambda, w}(t) - \sum_{n=-(m+1)}^{N-1} a_n(\lambda, w)t^n \right) dt + \sum_{n=-(m+1)}^{N-1} \frac{a_n(\lambda, w)}{s+n} + \int_1^\infty t^{s-1} \varphi_{\lambda, w}(t) dt.
$$

(52)

The last integral on the right hand side defines a function in $(s, w)$ that is holomorphic on

$$
\mathbb{C} \times \{ w \in \mathbb{C} \mid \text{Re} \ w > 0 \},
$$

the first integral defines a function holomorphic on

$$
\{ s \in \mathbb{C} \mid \text{Re} \ s > m + 1 - N \} \times \{ w \in \mathbb{C} \mid \text{Re} \ w > 0 \}.
$$

Thus, $(s, w) \mapsto \tilde{\phi}_{\lambda, w}(s)$ extends to a meromorphic function on

$$
\{ s \in \mathbb{C} \mid \text{Re} \ s > m + 1 - N \} \times \{ w \in \mathbb{C} \mid \text{Re} \ w > 0 \}.
$$

For any fixed $w \in \mathbb{C}$, $\text{Re} \ w > 0$, the poles of $\tilde{\phi}_{\lambda, w}$ are contained in

$$
\{ -(N - 1), \ldots, m + 1 \}.
$$

For $n \in \{ -(m + 1), \ldots, N - 1 \}$, the order of $s = -n$ as a pole is at most 1, its residue is $a_n(\lambda, w)$. Thus, if $\lambda \neq 1$ then the set of poles shrinks to

$$
\{ -(N - 1), \ldots, 0 \}.
$$

If $\lambda = 1$ then $a_{-(m+1)} = 1 \neq 0$, and $s = m + 1$ is indeed a pole.

This argumentation obviously applies to any $N$, and hence we find that $\tilde{\phi}_{\lambda, w}$ has a meromorphic continuation to all of $\mathbb{C}$ with all poles contained in $\mathbb{Z}_{\leq m+1}$, and each pole being simple. Since the Gamma function $\Gamma$ has simple poles at $-\mathbb{N}_0$, it follows that

$$
\Phi_m(s, \lambda, w) = \frac{1}{\Gamma(s)} \tilde{\phi}_{\lambda, w}(s)
$$

extends in $s$, for $\lambda \neq 1$, to an entire function, and, for $\lambda = 1$, to a meromorphic function on $\mathbb{C}$ with all poles contained in $\{ 1, \ldots, m + 1 \}$, all poles simple, and $m + 1$ being a pole. This completes the proof. \qed
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