RADIO-LOUD HIGH-REDSHIFT PROTOGALAXY CANDIDATES IN BOÖTES

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ABSTRACT

We used the Near Infrared Camera on Keck I to obtain Kσ-band images of four candidate high-redshift radio galaxies selected using optical and radio data in the NOAO Deep Wide-Field Survey in Boötes. Our targets have 1.4 GHz radio flux densities greater than 1 mJy, but are undetected in the optical to ≳24 Vega mag. Spectral energy distribution fitting suggests that three of these objects are at z > 3, with radio luminosities near the FR-I/FR-II break. The other has photometric redshift zphot = 1.2, but may in fact be at higher redshift. Two of the four objects exhibit diffuse morphologies in Kσ-band, suggesting that they are still in the process of forming.

Key words: galaxies: active – galaxies: high-redshift – galaxies: starburst – infrared: galaxies

Online-only material: color figures

1. INTRODUCTION

High-redshift radio galaxies (HzRGs) are rare objects in the cosmos, residing at the very brightest end of the radio luminosity function. Wide-area surveys have identified powerful radio galaxies out to very early cosmic epochs. However, since flux-limited surveys tend to probe increasingly luminous objects with increasing redshift, our knowledge of lower-luminosity HzRGs at high redshift is sparse (TN J0924-2201 has a 1.4 GHz flux density S1400 = 73 mJy, despite being at z = 5.19; van Breugel et al. 1999). We have sought to remedy this situation with deep, multifrequency radio maps of the Boötes field (de Vries et al. 2002; S. Croft et al. 2008, in preparation; Section 2). This field has substantial multwavelength supporting data.

Many HzRGs show diffuse rest-frame optical morphologies and other characteristics of “protogalaxies” (e.g., Penetucci et al. 2001); they appear to form when Lyman-break-galaxy-sized clumps merge and eventually evolve into large elliptical galaxies (van Breugel et al. 1998). Identifying lower-luminosity radio galaxies at early cosmic epochs can help our understanding of the nature of active galactic nucleus (AGN) hosts as a function of radio luminosity and can help constrain cosmic evolution of the radio galaxy population.

By applying a cut in radio spectral index, α (Sν ∝ να), we can preferentially select higher-redshift objects, since a combination of intrinsic evolutionary effects and an observational “k-correction” tend to cause α to become steeper (i.e., more negative) with increasing redshift (for details see, e.g., Krolik & Chen 1991). Indeed, 35% of sources with α < −1.3 and S1400 > 10 mJy observed by De Breuck et al. (2000) were found to have z > 3.

This paper presents Kσ-band images of radio sources with steeper than normal radio spectra and no detections in deep optical data to ≳24 mag (Vega) from the NOAO Deep Wide-Field Survey (NDWFS) Boötes field, in order to study the morphologies and properties of candidate HzRGs that are fainter than “typical” HzRGs.

Higdon et al. (2005) performed a complementary study of optically-invisible radio sources in the Boötes field. They used the Very Large Array (VLA) at 1.4 GHz to map a 0.5 deg² area to a limiting sensitivity (at field center) of ≈15 μJy beam⁻¹. Thirty six of their 377 radio sources (10%) were not visible in the NDWFS data, although 90% of these 36 had flux densities less than 1 mJy; in contrast (Section 3) we consider sources brighter than 1 mJy, and with relatively steep spectral indices, from our radio maps (Section 2) covering an order of magnitude larger area. Like us, however, Higdon et al. conclude that most of their optically-invisible objects are AGNs at relatively high redshifts. The region mapped by Higdon et al. has very little spatial overlap with our 325 MHz data, and none of their optically-invisible radio sources is present in our radio catalog.

Throughout this paper, we use Vega magnitudes, J2000 coordinates, and assume an Ωm = 0.24, ΩΛ = 0.76, H0 = 73 km s⁻¹ Mpc⁻¹ cosmology (Spergel et al. 2007).

2. RADIO DATA

Observations of a 4.9 deg² region of the Boötes field were made with the VLA A-array at 325 MHz during four runs from 2003 June–August. The CLEAN beam had a FWHM of 5″ and the limiting sensitivity was σ ≈ 150 μJy. The data analysis and reduction will be discussed in detail in an upcoming paper (S. Croft et al. 2008, in preparation). A catalog of radio sources (as opposed to components, i.e., a double-lobed radio galaxy counts as one source) was generated using the method described by de Vries et al. (2006).

We matched the 325 MHz catalog (S. Croft et al. 2008, in preparation) to the 1.4 GHz Westerbork Synthesis Radio Telescope (WSRT) dataset of de Vries et al. (2002), which has σ ≈ 28 μJy, an ellipsoidal CLEAN beam of 13″ × 27″, and an
area of 6.68 deg^2. The typical positional accuracy at 1.4 GHz for sources in the high signal-to-noise ratio (S/N) regime (which is the case for our sources since their fluxes are $> 1$ mJy) is 0\'0.44 (1\sigma). Six hundred and fifty three sources (in the 4.9 deg^2 where the two radio datasets overlap) were found to be in common and hence had measured radio spectral indices, $\alpha$. The properties of this matched catalog will be discussed by S. Croft et al. (2008, in preparation).

3. SAMPLE SELECTION

To select HzRG candidates, we combined deep radio data (Section 2) with optical data from the NDWFS (Jannuzi & Dey 1999), a deep ($B_W = 25.3, R = 24.1, I = 23.6, K \sim 19$ mag; $3\sigma, 5''$ diameter—Brodwin et al. 2006) optical/near-IR survey in Boötes. The magnitude limits for point sources are deeper, but HzRGs are likely to be resolved in the $\sim 1''$ seeing of NDWFS, so we used 4'' diameter optical apertures to determine whether or not a source is identified. We used the Brown et al. (2007) I-band selected catalog, which has improved photometry for faint objects compared to the public NDWFS DR3 Source Extractor (SExtractor) catalog. In computing photometric redshifts, we measured photometry from the NDWFS optical data in 4'' diameter apertures at each of the radio positions (Section 6). The NDWFS Boötes field is also covered by deep infrared, X-ray, and other observations (Section 5), although these data were not available to us at the time of our Near Infrared Camera (NIRC) observations and were not used as part of the selection.

As noted by Blumenthal & Miley (1979), Tielens et al. (1979), and others, radio sources with spectral indices that are steeper than normal (average $\alpha_{1400} \sim -0.7$) have a higher probability of being unidentified in optical surveys, a result that holds true for the combination of our radio data with the NDWFS optical data (Section 2) and steep spectral indices ($\alpha \lesssim 1$) have lower identification fractions than sources with "normal" spectral indices, suggesting that they are more distant and/or less luminous. This result will be discussed in more detail by S. Croft et al. (2008, in preparation).

Figure 1. Fraction of radio sources with $S_{1400} > 1$ mJy that are identified in the NDWFS $R$-band data, as a function of radio spectral index, $\alpha$. Radio sources with inverted ($\alpha \gtrsim 0.2$) and steep spectral indices ($\alpha \lesssim -1$) have lower identification fractions than sources with "normal" spectral indices, suggesting that they are more distant and/or less luminous. This result will be discussed in more detail by S. Croft et al. (2008, in preparation).

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Despite the efficiency of the method used for combining multiple radio source components in order to eliminate radio lobes from the catalog (Section 2), inevitably a few lobes may slip through. One such object turns out to be Source D; it was undetected in our NIRC observations and later turned out to be undetected in any of the IRAC bands. When we overplotted our VLA data, as well as data from the NRAO VLA Sky Survey (NVSS), on a wide-field Digitized Sky Survey (DSS) image, we found that this source is in fact a lobe of a large (4.5 diameter) FR-II that has a faint $R$-band identification in the DSS. Additionally, after the observations were taken, we replotted the radio data for Source B on the NDWFS images, including a lower contour than in our original plots, and noticed that the radio emission appeared extended (Figure 2), and so this object ought not to have been included in our original sample of 14 objects with unresolved radio emission and no NDWFS matches. Nevertheless, the radio emission appears to be associated with an IR-bright source which meets our criteria for optical faintness. Source D (the large FR-II) will not be discussed further here; the remaining four fields resulted in the detection of candidate HzRGs, as discussed in Section 6.

4. KECK NIRC/LRIS OBSERVATIONS

Imaging observations were made in $K_s$-band, using the NIRC (Matthews & Soifer 1994)—a 256 $\times$ 256 InSb array with 0\'15 pix$^{-1}$—on Keck I during two half nights on UT 2004 April 29.
and 30. A 16-position (4 × 4) dither, with dither spacing 3″ and 60 s per exposure (6 × 10 s), was used. Three fields were observed for four dithers (3840 s) and the remaining two fields observed for two dithers (1920 s), as summarized in Table 1. Seeing was 0′.5–0′.6 and the sky was clear. The data were reduced in the standard manner using DIMSUM10 in IRAF. Astrometry was bootstrapped from the NDWFS data; the small offsets between the $K_s$ and 4.5 μm positions (Table 1) suggest that our NIRC astrometry is accurate to $\sim$1″. Photometry was performed in a 3.5″ diameter aperture; the $K_s$ magnitudes are presented in Table 2, where we also present the measured NDWFS and ISS photometry (Section 6).

Two of the four HzRG candidates were observed spectroscopically. The slitmask was filled with secondary targets selected by their IR colors from the AGN “wedge” of Stern et al. (2005), along with a few HzRG candidates; these will be discussed in future papers. These observations used the Low Resolution Imaging Spectrometer (LRIS; Oke et al. 1995) on Keck I, with the D560 dichroic and a slit width of 1″. On the blue side, a 400 line mm$^{-1}$ grism blazed at 3400 Å was employed, giving 1.09 Å pixel and spectral resolution 8.1 Å. On the red side, a 400 line mm$^{-1}$ grating blazed at 8500 Å was used, giving 1.86 Å pixel and spectral resolution 7.3 Å. This setup gives spectral coverage of $\sim$3150–9400 Å, with some variation from slitlet to slitlet. Source C was observed on UT 2005 June 4 for 3600 s (2 × 1800 s) under clear skies, in 1″4 seeing. Source E was observed on UT 2005 June 6 for 4680 s (3 × 1560 s) in 0′8 seeing and with clear skies. The observations are summarized in Table 1.

The data were reduced in the standard manner using BOGUS11 in IRAF, and spectra extracted in a 1″5 wide aperture for many of the secondary targets. The primary targets were undetected in both the red and blue side two-dimensional spectra, and for these no extractions were performed.

5. OTHER DATA

5.1. X-ray

The Boötes field is covered by the Chandra X-ray observations of Murray et al. (2005). The resulting catalogs, published by Kenter et al. (2005) and Brand et al. (2006), reach a limiting full-band (0.5–7.0 keV) flux of $7.8 \times 10^{-15}$ erg s$^{-1}$ cm$^{-2}$ for on-axis sources with four counts. We checked for counterparts to our targets and found that Source E was detected12 (five counts in the full band, corresponding to 1.13 $\pm$ 0.71 $\times$ 10$^{-15}$ erg s$^{-1}$ cm$^{-2}$). This source is very likely to be a powerful high redshift and/or optically obscured quasar simply on the basis of its large X-ray-to-optical flux ratio (Brand et al. 2006).

Intriguingly, an object in the field of Source B is also detected in X-rays with five counts in the full band ($1.17 \pm 0.71 \times 10^{-15}$ erg s$^{-1}$ cm$^{-2}$). The X-ray position13 is close to the peak of the 325 MHz radio contours, but offset by (7′0) from the flux-weighted 1.4 GHz position. However, it appears that the X-ray and radio sources are two distinct objects; if they were the same (or if the radio source were associated with the object nearest the peak of the radio contours), the radio source would be extremely core dominated, which is not supported by the comparatively steep spectral index. Rather, it appears that the radio source is an extended, double-lobed source associated with the I-band dropout at the position given in Table 1, and that the X-ray source is associated with the optical object14 nearest (1′2) the X-ray position, and clearly visible in the NDWFS I-band (Figure 2) and NIRC $K_s$-band (Figure 7) data.

Both of the X-ray sources have too few counts to reliably determine a hardness ratio (three counts in the soft band and two in the hard band in both cases).

The other HzRG candidates are not detected in XBoötes, although they lie within the coverage area. Since the catalogs of Kenter et al. (2005) and Brand et al. (2006) contain only sources with four or more counts, it is possible that they may be X-ray sources, but be insufficiently bright for inclusion in the full catalog. None of our sources is within the coverage area of the more sensitive ($\sim$10$^{-16}$ erg s$^{-1}$ cm$^{-2}$), single-pointing Chandra X-ray observations of Wang et al. (2004).

5.2. Infrared

The zBoötes survey (Cool 2007) provides z-band data for part of the Boötes field. We checked for counterparts to our four HzRG candidates and found that three were undetected at the zBoötes flux limits (Table 2), while Source B was outside the zBoötes coverage area.

FLAMEX (Elston et al. 2006) provides J- and $K_s$-band catalogs for part of the Boötes field. The $K_s$-band catalogs are deeper than the $K_s$-band catalogs used for our selection, but our measured NIRC $K_s$-band fluxes (for the three sources in regions covered by FLAMEX) are below the FLAMEX $K_s$-band limit ($K_s \sim 19.3$ mag), so these sources are too faint to appear in the $K_s$-band FLAMEX catalog. The J-band FLAMEX limit ($J \sim 21$ mag) is insufficient to put anything but the weakest constraints on $J - K_s$ colors, and in any case, two of our sources are outside the FLAMEX coverage area. In future, refining

### Table 1

| Source name | I2000 IDa | 1.4 GHz R.A. | 1.4 GHz decl. | $K_s$ R.A. | $K_s$ decl. | 4.5 μm R.A. | 4.5 μm decl. | NIRC exp (s) | LRIS exp (s) |
|-------------|-----------|--------------|---------------|-----------|-----------|------------|-------------|-----------|-------------|
| Source A    | J142631+341557 | 14 26 31.75 | 34 15 57.57 | 14 26 31.8 | 34 15 57 | 14 26 31.75 | 34 15 58.5 | 3840 | 0 |
| Source B    | J142647+345851 | 14 26 47.87 | 34 58 51.0 | 14 26 47.8 | 34 58 53 | 14 26 47.78 | 34 58 53.3 | 1920 | 0 |
| Source C    | J142741+342324 | 14 27 41.84 | 34 23 24.7 | 14 27 41.9 | 34 23 25 | 14 27 41.85 | 34 23 26.3 | 3840 | 3600 |
| Source D    | J143036+340426 | 14 30 36.09 | 33 40 26.6 | No ID - radio source is a lobe | 1920 | 0 |
| Source E    | J143258+342055 | 14 32 58.44 | 34 20 55.4 | 34 20 56 | 34 20 58.44 | 34 20 56.0 | 1920 | 4680 |

Notes.
a J2000 identifiers are from the 1.4 GHz catalog of de Vries et al. (2002). All five sources listed correspond to 325 MHz counterparts without ambiguity. The de Vries et al. catalog is denoted [DMR2002] by SIMBAD.
b Detected in XBoötes X-ray data (Section 5.1).

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10 http://iraf.noao.edu/iraf/fp/contri/dimsumV2.
11 http://zwolfgang.kap.jpl.nasa.gov/~stern/homepage/bogus.html.
12 (R.A., decl.) = (14h32m58.49s, 34′20′′54.77″) $\pm$ 1.25″.
13 (R.A., decl.) = (14h26m48.41′, 34′58′′47.37″) $\pm$ 1.60″.
14 (R.A., decl.) = (14h26m48.45′, 34′58′′48.47″).
Figure 2. Postage stamps from NDWFS (BW RIK) and the IRAC Shallow Survey (3.6 and 4.5 μm) for the five sources in our Keck imaging sample. A 5′′ radius circle is plotted at the 1.4 GHz flux-weighted radio position (de Vries et al. 2002) in each frame, except K, where the 325 MHz radio map is overplotted. The lowest contour is at 0.4 mJy beam$^{-1}$; the remaining contours are evenly spaced in flux density up to the peak flux density in each map.

our selection criteria by requiring non-detections in FLAMEx will further increase our efficiency at finding HzRGs, since $K > 19$ corresponds to $z \gtrsim 2$ for sources on the $K - z$ relation (Figure 3; Section 6).

We checked the IRAC (Fazio et al. 2004) catalogs of the IRAC Shallow Survey (Eisenhardt et al. 2004) for infrared counterparts to our HzRG candidates. All were detected at greater than 2σ significance at 3.6 and 4.5 μm and one also at 5.8 μm. We report the catalog 5″ aperture magnitudes in Table 2 and use these measurements to assist with spectral energy distribution (SED) fitting (Section 6). Table 2 shows measured values, even in cases where the photometric uncertainties are large, since real measurements from the data reflect the relative weights given to the data points in the $\chi^2$ photometric redshift fitting. All objects are pointlike at IRAC resolution ($\sim 2″$).

We downloaded the post-BCD 24 μm mosaics from the Spitzer archival observations of Soifer et al. for the fields surrounding the four HzRG candidates and performed quick-
In order to obtain more stringent limits on the fluxes of our HzRG candidates (Table 2), which we use to make somewhat photometric redshift estimates. All magnitudes are on the Vega system.

![Figure 3](image_url)  
Figure 3. $K - z$ diagram of Willott et al. (2003) (circles represent radio galaxies with spectroscopic redshifts from the 3CRR, 6CE, 6C*, and 7CRS surveys). Our HzRG candidates are plotted as points with error bars, at the redshifts obtained from photometric redshift fitting. Vertical error bars represent the measured uncertainties in the $K_s$ band magnitudes (where we assume $K - K_s$ is negligible), and horizontal error bars represent the 68% confidence interval in redshift from HyperZ as tabulated in Table 3.

![Figure 4](image_url)  
Figure 4. $[3.6] - [4.5]$ versus $K_s - [3.6]$ diagram for our four HzRG candidates (magenta crosses with photometric error bars), compared to the colors of field sources in Boötes with at least a 5σ detection at NDWFS $K$-band (where we assume that $K - K_s$ is negligible), 3.6 µm and 4.5 µm (black points). Also plotted are the colors of non-evolving galaxy templates from Polletta et al. (2007), redshifted from $z = 0$ to $z = 7$ in steps of $z = 0.1$ (with redshifts marked in steps of $z = 1.0$). The red line is a 2 Gyr old elliptical template, the blue line is Arp 220, and the green line is a QSO2.

### 6. DISCUSSION

We have $K_s$ and IRAC detections and optical upper limits for our HzRG candidates (Table 2), which we use to make somewhat crude redshift estimates. $K_s - [3.6]$, which straddles the Balmer and/or 4000 Å breaks for the two $z_{phot} > 4$ sources from Table 3, provides quite tight redshift constraints. The position of our sources on a diagram of $[3.6] - [4.5]$ versus $K_s - [3.6]$ color (Figure 4) suggests that three are at $3 \lesssim z \lesssim 5$.

We also obtained photometric redshifts using the publicly available software HyperZ (Version 1.2, Bolzonella et al. 2000). In order to obtain more stringent limits on the fluxes of our objects than the 5′′ aperture limits used in the selection criteria, we remeasured photometry from the NDWFS $B_y R I$ images, in a 4″ aperture at the position of the radio centroid. Empirical photometric errors were derived using the method of Brown et al. (2007), a Monte Carlo technique which places apertures on random regions of sky. This provides more accurate estimates of uncertainties (verified by the addition of simulated galaxies to the real data) than methods which assume uncorrelated Poisson background noise and no source confusion. Although some of the resulting measurements have large uncertainties (Table 2), they help to constrain the resulting photometric redshifts better than using the published survey limiting magnitudes as upper limits. In cases where they are large, the quoted uncertainties provide a real measure of the variations in the background values due to neighboring objects. Since the uncertainties are used essentially as weights in the $\chi^2$ photometric fitting, large uncertainties downweight the effect of bands where the object

### Table 2

| ID    | $\alpha^{2000}_{\text{J}}$ | $S_{1400}$ (mJy) | $B_y$ | $R$ | $I$ | $z$ | $R_{\text{phot}}$ | $K_s$ | $[3.6]$ | $[4.5]$ | $[5.8]$ | $[8.0]$ |
|-------|--------------------------|------------------|-------|-----|----|-----|-------------------|-------|---------|---------|---------|---------|
| Source A | $-1.48$ | $1.08 \pm 0.04$ | $27.60 \pm 1.51$ | $25.06 \pm 0.64$ | $22.50 \pm 0.29$ | $>23.1$ | $>20.8$ | $20.94 \pm 0.10$ | $18.12 \pm 0.14$ | $17.34 \pm 0.14$ | $15.84 \pm 0.31$ | $15.98 \pm 0.57$ |
| Source B | $-0.89$ | $8.93 \pm 0.04$ | $28.25 \pm 1.76$ | $25.91 \pm 0.91$ | $27.87 \pm 3.36$ | $\cdots$ | $19.63 \pm 0.05$ | $18.03 \pm 0.13$ | $17.58 \pm 0.18$ | $15.75 \pm 0.47$ |
| Source C | $-0.98$ | $3.61 \pm 0.03$ | $26.70 \pm 0.73$ | $25.32 \pm 1.12$ | $24.40 \pm 0.81$ | $>22.9$ | $\cdots$ | $21.38 \pm 0.13$ | $19.23 \pm 0.37$ | $18.51 \pm 0.40$ | $17.27 \pm 1.13$ | $16.02 \pm 0.60$ |
| Source E | $-0.87$ | $16.27 \pm 0.03$ | $26.37 \pm 0.76$ | $25.73 \pm 1.24$ | $24.48 \pm 1.10$ | $>23.2$ | $>20.6$ | $20.99 \pm 0.13$ | $18.45 \pm 0.19$ | $17.88 \pm 0.23$ | $17.08 \pm 0.95$ | $16.42 \pm 0.85$ |

**Notes.** We report measured photometric uncertainties, even when these are large (e.g. 6.05 mag for Source B at 5.8 µm), since these were the values used in photometric redshift fitting (Section 6). All magnitudes are on the Vega system.

* NDWFS (Section 3).
* zBoötes (Section 5.2).
* FLAMEX (Section 5.2).
* NIRC (Section 4).
* IRAC Shallow Survey (Section 5.2).
flux is poorly constrained, as opposed to assigning these bands too high a weight by setting the fluxes to zero and obtaining the uncertainty from the survey flux limit. However, even when we re-run the photometric redshift fitting using the survey limits for those sources with uncertainties larger than 1 mag, we find that there is essentially no change in the best-fit redshift (the largest change is for Source E, whose best-fitting template changes from Arp220 to a 2 Gyr old elliptical, and whose redshift changes by just 0.04), although the confidence intervals become marginally less well constrained. For the z- and J-bands, where we only have access to catalogs and not to the original images (in order to re-measure photometry at the radio positions), we set the fluxes to zero and derive the associated uncertainties from the published survey limiting magnitudes. We find that whether or not the z- and J-band data points are included has essentially no effect on the best-fitting redshifts and confidence intervals (due to the large uncertainties in these bands).

The J-band point for Source A seems rather high (Figure 5); inspection of the NDWFS J-band postage stamp in Figure 2 suggests that this is probably due to the unrelated saturated star to the southwest. In fact, even though Source A fulfils our initial selection criterion (no source in the Brown et al. 2007 catalog closer than 4″), the aperture magnitude measured at the radio position (Table 2) is brighter than the 3σ survey flux limits quoted in Section 3, strongly suggesting that light from a nearby object is contributing to the measured aperture magnitude (the NDWFS J-band point-spread function (PSF) has wings which can cause problems for aperture photometry at surprisingly large radii). If the photometric redshift is correct, there could also be some contribution to the flux in this band from Lyα at z ≈ 5, although an I = 22.5 source ought to be visible in the postage stamps (Figure 2). Excluding the J-band point from the SED fitting for this object has no effect on the best-fitting redshift and little effect on the confidence intervals.

Instead of using the templates provided with HyperZ, we used a selection of templates (QSO2, Arp220, Mrk231, EII2, EII5, and EII13 templates; the latter three are 2, 5, and 13 Gyr old ellipticals) from Polletta et al. (2007), which do a better job of reproducing IR fluxes. The effects of dust extinction were simulated using the prescriptions of Calzetti et al. (2000), with extinction $A_V$ allowed to vary between 0.0 and 2.0 mag (although some of the templates already incorporate the effects of some intrinsic dust extinction). $B$-band absolute magnitude, $M_B$, was constrained to be $-28.8 \leq M_B \leq -19$. Because of the faintness of these objects in the optical, and the corresponding large uncertainties for the optical data points compared to the IR, the range of plausible redshifts is rather large (Table 3). However, we can still estimate the best-fitting redshifts, $z_{\text{phot}}$. The template SEDs for the best-fitting models are shown in Figure 5.

Three of the four HzRG candidates have $z_{\text{phot}} > 3$, and the fourth is rather poorly constrained (due to low S/N) to $z_{\text{phot}} = 1.2$ (Table 3). We report precise values for $z_{\text{phot}}$ in order that the reader may know what values were used for the calculation of the rest-frame radio luminosities and other such properties. However, as can be seen from the confidence limits in redshift reported in Table 3, we cannot constrain redshifts for these objects very accurately using SED fitting, given the large number of bands with non-detections or large photometric errors. Source B was not observed spectroscopically, so we cannot say whether or not the emission lines seen in the template (Figure 5) are really present in its spectrum. In the event of galaxy-wide dust extinction these spectral lines would be less prominent, but this ought not to affect our determination of $z_{\text{phot}}$ too much.

We can also use the 4.5 μm flux densities alone as an alternative to the $K - z$ relation. Seymour et al. (2007) studied a sample of 69 HzRGs with spectroscopic redshift 1 < $z$ < 5.2 and a wide range of radio luminosities. They showed that the rest-frame $H$-band light for these objects is well described by models of elliptical galaxies with formation redshift $z_f = 10$ and masses between $10^{11}$ and $10^{12}$ $M_\odot$. In Figure 6, we plot the measured values of the 4.5 μm flux density from Seymour et al. against redshift, along with tracks representing Bruzual & Charlot (1993) models with $z_f = 10$ and solar metallicity, normalized to $10^{11}$ and $10^{12}$ $M_\odot$. Plotting our HzRG candidates in this figure at their best-fit photometric redshifts, these redshifts seem reasonable, although the lower bounds on the redshift (assuming our objects have masses > $10^{11}$ $M_\odot$) are still quite weak ($z \gtrsim 2$). Source C, with $z_{\text{phot}} = 1.21$, falls below the $10^{11}$ $M_\odot$ line, suggesting that it is either unusually low mass, or that the SED fitting underestimates the true redshift (possibly certain given the wide confidence interval and low S/N for this source). Source C is the only outlier on the $K - z$ plot (Figure 3), again suggesting it is either very underluminous, or is in fact at higher redshift than suggested by the SED fits.

The NIRC images (Figure 7) provide much higher resolution than IRAC, and we are able to see that two of the detected sources have somewhat extended, “fuzzy” morphologies (Figure 8). Source E is somewhat more compact, but appears to have a faint trail of material extending to the northeast. Source B seems to be quite compact and relatively bright. The extended sources have $K_s$-band morphologies suggestive of the optical morphologies seen in HzRGs such as B2 0902+34 (van Breugel et al. 1998), which hints that these objects may be galaxies still in the process of formation (again arguing for a higher redshift for Source C than that obtained from SED fitting). The compactness of the observed radio morphologies (unresolved in our maps, except for Source B) suggests that we are probably seeing these radio galaxies relatively soon after the AGN

### Table 3

| ID       | $z_{\text{phot}}$ | $z$ (68% confidence) | $z$ (99% confidence) | Template  | $A_V$ | $M(B_{\text{rest}})$ | Rest-frame 1.4 GHz $L$ ($10^{26}$ W Hz$^{-1}$) |
|----------|-------------------|----------------------|----------------------|-----------|-------|----------------------|-----------------------------------------------|
| Source A | 4.97              | 4.92–4.98            | 3.94–5.69            | Arp220    | 0.00  | −24.67              | 6.94                                          |
| Source B | 3.76              | 3.68–3.81            | 1.77–4.08            | QSO2      | 0.30  | −24.32              | 10.5                                          |
| Source C | 1.21              | 0.00–2.76            | 0.00–6.99            | Mrk231    | 0.00  | −19.05              | 0.298                                         |
| Source E | 4.65              | 4.21–4.99            | 0.00–6.99            | Arp220    | 0.00  | −24.07              | 30.2                                          |

**Note.** Precise values of $z_{\text{phot}}$ are reported, even when the associated uncertainties are large, so that the reader may know the exact redshifts used to compute rest-frame properties of the sources.
Figure 5. SEDs of the HzRG candidates. The best-fitting templates are shown plotted at the redshift which gave the lowest $\chi^2_{\nu}$ in SED fitting. We also note the best-fit redshift in each panel; note, however, that the redshifts are not very tightly constrained (Table 3).

(A color version of this figure is available in the online journal)

Figure 6. 4.5 $\mu$m flux density versus redshift for HzRGs with spectroscopic redshifts from Seymour et al. (2007) (circles) and our HzRG candidates (crosses, plotted at the best-fit photometric redshift with 68% confidence intervals shown). The dashed and dotted lines are evolutionary tracks for elliptical galaxies with formation redshift $z_f = 10$ and masses $10^{11}$ and $10^{12} M_\odot$, respectively. (A color version of this figure is available in the online journal)

Radio galaxies are among the most massive objects at all redshifts (Seymour et al. 2007), but there is some evidence for a dependence of host optical luminosity on radio power (e.g., De Breuck et al. 2002; McLure et al. 2004). It is plausible that our objects are less powerful in the radio and hence might be expected to be fainter than the $K - z$ relation (Figure 3); they might, in fact, be still forming the bulk of their stars (as suggested by their morphologies) but the AGN has already switched on. Source B, which is more compact in the near-IR, and has more extended radio emission (suggesting that the AGN activity is less recent than in the other three sources), is slightly brighter than the $K - z$ relation, again as might be expected in this scenario.

From the measured radio fluxes and spectral indices, and the best-fit photometric redshifts, we calculate rest-frame 151 MHz...
luminosities (for the $z_{\text{phot}} > 3$ sources) of $\sim 1 \times 10^{27}$ W Hz$^{-1}$ sr$^{-1}$, which implies (Willott et al. 1999; McCarthy 1993) Ly$\alpha$ luminosities $\sim 2 \times 10^{41}$ erg s$^{-1}$. The corresponding Ly$\alpha$ fluxes are $\sim 6 \times 10^{-17}$ erg s$^{-1}$ cm$^{-2}$, essentially independent of redshift (due to similar effects of distance on the observed Ly$\alpha$ and radio fluxes, and hence the determination of rest-frame luminosities). A typical HzRG Ly$\alpha$ line (rest-frame FWHM $\sim 5$ Å) would be marginally unresolved in our spectra at $z = 0$, but beyond $z \sim 0.6$, observed flux densities would start to decrease due to broadening of the Ly$\alpha$ line in the observed frame. At $z \sim 5$, the expected Ly$\alpha$ flux density is $2 \times 10^{-18}$ erg s$^{-1}$ cm$^{-2}$ Å$^{-1}$. Since the noise in our spectra has $\sigma \sim 4 \times 10^{-19}$ erg s$^{-1}$ cm$^{-2}$ Å$^{-1}$, the predicted S/N is rather low ($\sim 5$), so it is perhaps unsurprising that we do not observe emission lines in our spectra, especially given the uncertainties ($\pm 1$ dex) in the emission line—radio correlation. If the line happens to fall in a region of the spectrum where noise is higher (at the position of a subtracted telluric emission line, for example), the S/N would be lower still. Dust, if present in the source, would also reduce the predicted flux densities and hence the S/N. Source A has the steepest radio spectrum of our candidates, so is most similar to classical HzRGs, and remains a good candidate for future spectroscopy. As discussed above, Source B may have prominent emission lines, and spectroscopic observations would be worthwhile.

The photometric redshifts and “typical” HzRG morphologies support the interpretation of these objects as HzRGs (and their deduced radio luminosities and $B$-band absolute magnitudes are AGN-like, as expected). At the fainter radio flux levels we probe, we seem to be seeing the more “normal” members of the HzRG population that perhaps make up the bulk of the luminosity function at these redshifts. Indeed, comparing to the results of Ledlow & Owen (1996), we find that our three $z > 3$ galaxies (with rest-frame $R$-band magnitudes $M_R \sim -25.5$ mag, as computed from the photometric redshifts in Table 3) ought to lie on the FR-I/FR-II break ($\sim 10^{27}$ W Hz$^{-1}$ at these optical absolute magnitudes).

As a sanity check, we obtain a rough estimate of the space density of these objects by calculating the comoving volume covered by the 4.9 deg$^2$ combined radio dataset, in the redshift range 1.21–4.97 (the range of photometric redshifts in Table 3). In our chosen cosmology, this volume is $2.3 \times 10^{8}$ Mpc$^3$. Since we do not claim to have exhaustively studied the radio-loud optically-faint galaxy population in this field within this redshift range, the resulting space density for our four HzRG candidates of $2 \times 10^{-8}$ Mpc$^{-3}$ should be taken as a lower limit. This compares favorably with the lower limit on radio galaxies of similar radio power ($P_{1400} > 10^{27}$ W Hz$^{-1}$) but steeper spectral index ($\alpha \lesssim -1.3$) obtained by De Breuck et al. (2006), $1.2 \times 10^{-9}$ Mpc$^{-3}$ at $3 < z < 4$. Since only one out of four of our HzRG candidates has a spectral index as steep as this (and since both estimates are lower limits), the space density is consistent with that measured by De Breuck et al.
Deeper optical and/or IR observations (particularly $J$, $H$ and/or mid-IR) would provide better constraints on our photometric redshifts. As discussed above, our objects may be only marginally too faint to be detected in our LRIS observations, so it is possible that deeper (red-sensitive) 10 m spectroscopy could in future provide spectroscopic redshifts. Mid-IR observations would also be of interest in constraining star-formation rates, and the presence and temperature of any dust in these objects.

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