High-throughput computational screening of nanoporous adsorbents for CO₂ capture from natural gas†

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With the growth of natural gas as an energy source, upgrading CO₂-contaminated supplies has become increasingly important. Here we develop a single metric that captures how well an adsorbent performs the separation of CH₄ and CO₂, and we then use this metric to computationally screen tens of thousands of all-silica zeolites. We show that the most important predictors of separation performance are the CO₂ heat of adsorption \(Q_{st,CO₂}\) and the CO₂ saturation loading capacity. We find that a higher-performing material results when the absolute value of the CH₄ heat of adsorption \(Q_{st,CH₄}\) is decreased independently of \(Q_{st,CO₂}\), but a correlation that exists between \(Q_{st,CH₄}\) and \(Q_{st,CO₂}\) in all-silica zeolites leads to incongruity between the objectives of optimizing \(Q_{st,CO₂}\) and minimizing \(Q_{st,CH₄}\), rendering \(Q_{st,CH₄}\) nonpredictive of separation performance. We also conduct a large-scale analysis of ideal adsorbed solution theory (IAST) by comparing results obtained using directly-generated mixture isotherms to those obtained using IAST; IAST appears adequate for the purposes of establishing performance trends and structure-property relationships in a high-throughput manner, but it must be tested for validity when analyzing individual adsorbents in detail since it can produce significant errors for materials in which there is site segregation of the adsorbate species.

1. Introduction

Natural gas, composed largely of \((\text{CH}_4)\), is a large and growing part of the domestic and global energy portfolio, with total domestic production having grown by 35% from 2005 to 2013 and now representing 28% of total American energy consumption.¹ Although this growth occurred for largely economic reasons as a response to the development of shale gas resources,¹ since the carbon dioxide \((\text{CO}_2)\) emissions per unit of electricity are about half as large from natural gas power plants as from coal power plants, a shift from coal to natural gas may be an effective way to reduce anthropogenic greenhouse gas emissions during the transitory period toward a sustainable energy economy,² though the practical achievement of these environmental benefits requires adequate...
controls of fugitive gas emissions. Additionally, since natural gas power plants have lower cycling times and costs than coal power plants, they can adjust more quickly to fluctuations in grid electricity demands: an advantage of increasing importance given the intermittency of many renewable energy sources such as wind and solar power.

Prior to use, the natural gas must be treated to remove contaminants including acid gases such as CO$_2$ and hydrogen sulfide, inert gases such as nitrogen (N$_2$) and helium, and other species such as water and heavier hydrocarbons. CO$_2$ is one of the most common impurities, and in addition to its presence decreasing the natural gas’s heating value, it can cause problems during transportation by corroding pipelines or forming solids in cryogenic tanker trucks, necessitating purity specifications like $<2$ mol\% CO$_2$ for a natural gas pipeline.

Porous solids can be used for the separation of CH$_4$ and CO$_2$ via the processes of pressure-swing adsorption (PSA), temperature-swing adsorption (TSA), or a hybrid process (PTSA), in all of which the adsorption columns undergo a cycle of adsorption followed by a regenerating desorption cycle at conditions of decreased pressure and/or increased temperature. In order to find the optimal adsorbents for a particular separation, metrics are necessary to quantify an adsorbent’s separation performance. Two of the most widely used metrics are the adsorption selectivity, defined as the ratio of the loading mole fractions divided by the ratio of the gas phase mole fractions, and the working capacity for each of the adsorbates, defined as the difference in loading of each adsorbate at adsorption and desorption conditions. These metrics are easily calculated from mixture adsorption isotherms, and many adsorbent screening projects have examined these metrics individually to evaluate an adsorbent at given adsorption and desorption conditions. Other metrics have been proposed as being more appropriate for the comparison of adsorbents. Rege and Yang and Wiersum et al. suggested multiplying the selectivities and working capacities in different manners to obtain metrics they called the sorbent selection parameter (S) and the adsorbent performance indicator (API). As we will show, these metrics can lead to undesirable results since they directly include the selectivity, which is unbounded and can approach infinity. Some studies have created metrics that aim to provide a more direct representation of the process economics. For example, in the CO$_2$/N$_2$ gas separation being investigated for carbon capture and storage (CCS) from the flue gas of fossil-fuel powered power plants, the economic driver is the energy cost of the separation, which has been estimated to parasitically reduce the net output of power plants by 30\%. This parasitic energy has been simply estimated as a sum of the thermal energy required to desorb the adsorbate and of the compression work by Lin et al. and Maring and Webley has developed a simplified process model to estimate the power required for the vacuum pump. Of course, a full process modeling analysis provides the most realistic cost estimates, but such detailed process modeling can be computationally expensive and thus unfeasible for a high-throughput screening study. Additionally, such process modeling requires expertise beyond what many researchers studying adsorption science possess.

In this contribution, we first develop a single simple metric that is representative of the economic drivers behind the CH$_4$/CO$_2$ gas separation required for upgrading various natural gas sources to pipeline quality, which we call the separation performance parameter (SPP). We then use this metric to explore the separation performance and structure–property relationships of tens of thousands of all-silica zeolites taken from the International Zeolite Association (IZA) database and the predicted crystallography open database (PCOD) of hypothetical zeolites, using a high-performance GPU grand-canonical Monte Carlo (GCMC) code to directly-generate mixture isotherms. Finally, we take the opportunity afforded by the large amount of data collected in this study to evaluate the accuracy of ideal adsorbed solution theory (IAST), a prominent technique to generate mixture adsorption isotherms. Details of the methods used in this work can be found in the ESL.

We considered three different sources of natural gas, all at a temperature of 300 K and assumed to be a binary mixture of CH$_4$ and CO$_2$: (1) landfill fill gas (LFG), produced by microorganisms acting upon municipal wastes, assumed to contain 60 mol\% CH$_4$ at a total pressure of 1 bar, (2) associated petroleum gas (APG), found in oil reservoirs, assumed to contain 90 mol\% CH$_4$ at a total pressure of 7 bar, and (3) non-associated gas (NAG), found apart from oil, assumed to contain 90 mol\% CH$_4$ at a total pressure of 70 bar. These pressures and temperatures were used as the adsorption conditions. As different materials perform optimally at different conditions, we evaluated each material at desorption conditions that optimize its SPP to allow for fair comparison.

2. Metric development

We chose to model the natural gas upgrading process as one in which CO$_2$ is selectively adsorbed for several reasons. As CO$_2$ tends to adsorb onto zeolites more strongly than CH$_4$ due to its quadrupole moment, more zeolites will be capable of performing this separation, with about 90\% of the hypothetical zeolites having a CO$_2$ selectivity greater than 1 at conditions of 1 bar total pressure, 300 K, and 90 mol\% CH$_4$. Furthermore, in making the raffinate the high-value stream, there is no energetic cost to performing PSA down to
atmospheric pressure since there is no need to compress the product, unlike in CCS and other adsorption processes where the adsorbed gas is the desired product. Finally, in adsorbing the minor component in the processes we analyzed, cycle times and energetic costs are expected to be less.

Upgrading natural gas requires the examination of multiple costs to evaluate the economics of different adsorbents. The fixed costs vary with the size of the adsorption column and thus the amount of adsorbent required for the separation. The variable costs include the energy required for operating the adsorption column. A high fractional CH₄ recovery is also desirable to reduce the variable cost of the feed gas as well as to reduce the loss of CH₄ to the atmosphere with the waste CO₂. If equilibrium adsorption and desorption are assumed, i.e., the mass transfer zone length approaches 0, these metrics can be individually calculated given mixture adsorption isotherms, as shown below. They can also be grouped into a single objective function to form a multiple-input single-output optimization problem, with the inputs being the desorption pressure (PSA), desorption temperature (TSA), or both (PTSA), and the output being the objective function value (OFV). By varying the inputs to minimize the OFV, one finds the optimal desorption conditions for a given adsorbent.28 The better-performing adsorbents will have a lower OFV, which leads us to term the inverse of the OFV the separation performance parameter (SPP), which we use to evaluate the screening results.

It would seem that the objective function that most appropriately captures the process economics would be a summation of the capital and operating expenditures of the separations plant. However, our goal of creating a simple and non ephemeral metric (the value of which is not dependent on present market conditions) precludes such a calculation. Instead, we will assume that the mass of adsorbent, energy, and fractional CH₄ recovery are equally important economic drivers, and we will then go on to show that this assumption results in a metric that closely agrees with the economic results of a more detailed process simulation study. Thus we use as a definition of the OFV and SPP:

\[
\text{OFV} = \frac{1}{\text{SPP}} = \left( \frac{M_{\text{ads}}}{M_{\text{CH}_4,\text{eff}}} \right) \times \frac{E}{M_{\text{CH}_4,\text{eff}}} \left( \frac{M_{\text{CH}_4,\text{eff}}}{M_{\text{CH}_4,\text{feed}}} \right)
\]

where \(M_{\text{ads}}\) is the mass of the adsorbent, \(M_{\text{CH}_4}\) is the moles of species \(i\) in stream \(k\), \(E\) is the total energy required for the separation, “feed” is the adsorption column inlet stream, and “raff” (raffinate) is the adsorption column outlet stream consisting of the gas that did not adsorb which contains the recovered CH₄. Note that here and throughout this paper, we use units of mass for the adsorbent and units of moles for the adsorbates. The three terms in eqn (1) are given below, with more details provided in the ESI.†

The first term in eqn (1) is the mass of adsorbent required per mole of CH₄ captured in one batch adsorption–desorption cycle:

\[
\frac{M_{\text{ads}}}{M_{\text{CH}_4,\text{eff}}} = \frac{y_{\text{CH}_4,\text{raff}} - y_{\text{CH}_4,\text{feed}}}{y_{\text{CH}_4,\text{feed}} \left( \Delta q_{\text{CH}_4} + \Delta q_{\text{CO}_2} \right) - \Delta q_{\text{CH}_4}}
\]

where \(y_{\text{CH}_4,\text{raff}}\) is the mole fraction of species \(i\) in stream \(k\) and \(\Delta q_{\text{CH}_4}\) is the working capacity for species \(i\). Since this is not the way that adsorption columns are sized in practice, \(M_{\text{ads}}/M_{\text{CH}_4,\text{eff}}\) can alternatively be thought of as a proxy for the cycle time of a batch when comparing equally-sized columns. Eqn (2) may be further understood in the context of two extremes. A perfectly selective adsorbent will have \(\Delta q_{\text{CH}_4} = 0\), and if \(y_{\text{CH}_4,\text{raff}}\) is set to 1, then \(M_{\text{ads}}/M_{\text{CH}_4,\text{eff}}\) is simply equal to \(\Delta q_{\text{CO}_2}/\Delta q_{\text{CH}_4}\).

At the other extreme, as a material’s ratio of \(\Delta q_{\text{CH}_4}/\Delta q_{\text{CO}_2}\) increases, approaching the ratio of \(y_{\text{CH}_4,\text{feed}}/y_{\text{CO}_2,\text{feed}}\) less and less separation will be performed and \(M_{\text{ads}}/M_{\text{CH}_4,\text{eff}}\) will approach infinity. Materials with a working capacity ratio exceeding this limit will have a required mass of adsorbent that is negative, and cannot perform the given separation.

If \(y_{\text{CH}_4,\text{raff}}\) is set to 1, one obtains the mass of adsorbent required to capture a mole of CH₄ in an adsorption column operated until the column is completely saturated. Alternatively, one may set \(y_{\text{CH}_4,\text{raff}}\) to the purity required for a specific application, which corresponds to an adsorption column with a mass transfer zone length approaching 0 that is operated beyond initial breakthrough in order to obtain more product which on average still meets its specifications. The latter option was chosen for this study, and \(y_{\text{CH}_4,\text{raff}}\) was set to 98 mol% CH₄.

The second term in eqn (1) is the energy required per mole of CH₄ captured:

\[
\frac{E}{M_{\text{CH}_4,\text{eff}}} = \frac{M_{\text{ads}}}{M_{\text{CH}_4,\text{eff}}} \left[ C_p \left( T_{\text{ads}} - T_{\text{des}} \right) + \Delta q_{\text{CH}_4} Q_{\text{ads}} + \Delta q_{\text{CO}_2} Q_{\text{des}} \right] + \frac{W_{\text{vac}}}{M_{\text{CH}_4,\text{eff}}}
\]

where \(C_p\) is the heat capacity of the adsorbent (assumed to be 0.75 kJ kg⁻¹ K⁻¹ for all zeolites),\( T_{\text{ads}}\) is the desorption temperature, \(T_{\text{des}}\) is the adsorption temperature, \(Q_{\text{ads}}\) is the isosteric heat of adsorption at zero loading of species \(i\), and \(W_{\text{vac}}\) is the work required to pull vacuum. Here and throughout this work, \(Q_{\text{des}}\) refers to the negative value of the enthalpy of adsorption such that positive values are shown. The energy term in eqn (3) consists of the sensible energy to raise the temperature of the bed during a temperature swing, the energy to desorb adsorbates, and the work to pull vacuum on
the adsorption column. We approximate the work required to pull vacuum below 1 bar as the energy required to compress the desorbed vapor from its vacuum pressure to atmospheric pressure using the formula for isentropic compression:14,15,31

\[
\frac{W}{M_{\text{CH}_4,\text{eff}}^{\text{SPP}}} = \begin{cases} 
\frac{1}{\eta} R T_{\text{des}} \gamma \left( \frac{P_{\text{atm}}}{P_{\text{des}}} \right)^{\gamma - 1} \left( \Delta q_{\text{CH}_4} + q_{\text{CO}_2} \right) M_{\text{atm}} M_{\text{CH}_4,\text{eff}}^{\text{SPP}}, & \text{if } 0.1 \text{ bar} < P_{\text{des}} < 1 \text{ bar} \\
0, & \text{if } 1 \text{ bar} < P_{\text{des}} 
\end{cases}
\]

(4)

where \( \eta \) is the vacuum pump efficiency (assumed to be 75%), \( R \) is the gas constant, \( \gamma \) is the heat capacity ratio (assumed to be 1.3 for both CH4 and CO2), \( P_{\text{atm}} \) is atmospheric pressure (approximated as 1 bar), and \( P_{\text{des}} \) is the desorption pressure.32 Eqn (4) was considered valid down to a desorption pressure of 0.1 bar; at lower pressures, the equation may fail as the efficiency falls below 75%.31

The third term in eqn (1) is the fractional CH4 recovery:

\[
\frac{M_{\text{CH}_4,\text{eff}}}{M_{\text{CH}_4,\text{feed}}} = \frac{\gamma_{\text{CH}_4,\text{feed}} \left( \Delta q_{\text{CH}_4} + \Delta q_{\text{CO}_2} \right)}{\gamma_{\text{CH}_4,\text{eff}} \left( \Delta m_{\text{CH}_4} + \Delta m_{\text{CO}_2} \right) - \Delta q_{\text{CH}_4}}
\]

(5)

Using eqn (2)-(5) we can calculate the OFV and SPP in eqn (1) at a given desorption pressure and desorption temperature. The components of the objective function given by these equations can be individually evaluated, but they cannot necessarily be individually used as the full objective function to find the optimal desorption conditions because this can result in desorption conditions that are at the most extreme values allowed. For example, if one were to attempt to minimize the mass of absorbent required for a PSA process, one would arrive at a desorption pressure of 0.1 bar for every material as there would be no energetic penalty.

We have placed the computer code we developed for the screening online with the Open Science Framework, along with the pure-component and mixture isotherms and the screening results.33 The code is capable of computing the SPP at optimized or set desorption conditions either by using input mixture isotherms or by using input pure-component isotherms to calculate mixture isotherms with IAST.

Separation performance parameter validity

At this point it is important to mention that a requirement imposed during the development of the SPP was that it be practical to compute the metric for thousands of materials, which precludes a full process design for all zeolites. We intended for the SPP to be representative of the most important economic drivers behind the CH4/CO2 gas separation. However, the particular form of the SPP (a product of factors) instead of the conventional sum of capital and operating costs assumes that all costs scale with the amount of adsorbent. In some important cases a significant fraction of the capital costs is independent of the amount of material, and if these capital costs dominate, the amount of material can become irrelevant. In addition, our study used isothermal working capacities to calculate the SPP for PSA processes whereas industrial adsorption columns operate under adiabatic conditions, which leads to lower working capacities.14 It is therefore important to compare the SPP with the separation costs calculated by a more detailed process engineering analysis which used adiabatic working capacities. First et al.16 conducted just such a study. From a database of 199 IZA zeolites, the authors selected 86 that topological analysis suggested would be viable for the CH4/CO2 separation, and adsorption isotherms indicated that 22 of these were highly selective for CO2. These 22 zeolites then underwent PSA process modeling and optimization to minimize the separation cost. Of the processes First et al.16 analyzed, the most similar to one of the three processes we analyzed was one which had adsorption pressures of 3 to 5 bar, desorption pressures of approximately 0.1 bar, and 90 mol% CH4; conditions quite similar to the APG process. The authors found that 8 zeolites could most feasibly perform the separation (ABW, AEN, AHT, APC, BIK, JBW, MON, and WEI), with the separation cost being practically identical among them. Of these top 8 performing zeolites, the SPP identified 4 as the top 4 materials of the 110 IZA zeolites we investigated, and the remaining 4 as being within the top 20 (Table S2†). These results indicate that the SPP is a reasonable indicator of separation performance.

To further justify the need to use the SPP, we can compare it to alternative metrics that have been given earlier, such as the sorbent selection parameter (S) of Rege and Yang:17

\[
S = \frac{\sigma_{\text{CO}_2,\text{CH}_4,\text{ads}}^2}{\sigma_{\text{CO}_2,\text{CH}_4,\text{des}}} \frac{\Delta m_{\text{CO}_2}}{\Delta m_{\text{CH}_4}}
\]

(6)

where \( \sigma_{ij} = \frac{q_i}{Y_i,\text{feed}/Y_j,\text{feed}} \) is the selectivity for species i over species j at adsorption or desorption conditions. The sorbent selection parameter seemingly has the advantage of being simpler than the SPP in eqn (1), though the only additional data required by eqn (1) are the heats of adsorption and an estimate for the specific heat of the adsorbent (with the latter being unnecessary for the evaluation of a PSA process). Indeed, we found that the sorbent selection parameter compared as favorably with the results of First et al.16 as the SPP, with the 8 top performing zeolites all being within the top 20 zeolites as ranked by the sorbent selection parameter (2 were tied for first with 11 zeolites that had infinite selectivity for CO2, and 6 were ranked within 12–20), so a justification for the value of the SPP is wanted.

Fig. 1 shows the relationship between the SPP developed in this work and the sorbent selection parameter for the APG process carried out with PSA at a set desorption pressure, which is representative of the results for the other two
Fig. 1 A correlation between the SPP and the sorbent selection parameter for the APG process carried out with PSA at set desorption conditions of 1 bar and 300 K. Hypothetical zeolites are shown as black dots, IZA zeolites are shown as blue diamonds, and the hypothetical zeolites in Table 1 are shown as green circles.

processes as well. A close correlation between the two metrics can be observed, though the metrics begin to differ for the top-performing materials, with the sorbent selection parameter being more variable than the SPP due to the selectivity being unbound. We examine a subset of materials selected for their deviating SPP and sorbent selection parameter in Table 1. Here, the undesirability of having a performance metric dependent on the unbound selectivity is demonstrated as all the materials in Table 1 have sufficiently low Δq_CH4 values such that additional decreases in Δq_CH4 should have little effect on separation performance, yet the sorbent selection parameter continues to increase as Δq_CH4 is negligibly decreased despite much more significant decreases in Δq_CO2; for example, compare PCOD8294501 and PCOD8310046 in Table 1: both have Δq_CH4 two orders of magnitude lower than Δq_CO2 and thus effectively zero, yet the sorbent selection parameter of PCOD8294501 is greater than that of PCOD8310046 despite the latter’s much larger Δq_CO2. Furthermore, for materials which do not load any CH4 at all, the selectivity for CO2 goes to infinity and yields sorbent selection parameter that are uncomparable. Overall, the use of the sorbent selection parameter can thus mislead scientific efforts focused on finding the highest-performing materials.14

In addition, the sorbent selection parameter cannot be used to find the optimal desorption conditions for a given adsorbent,14 as we found that the ratio of working capacities was almost always more sensitive to a change in desorption conditions than the selectivity at desorption conditions (the selectivity at adsorption conditions does not change as a function of desorption conditions). For example, with PSA the sorbent selection parameter was almost always maximized at either the lowest or the highest allowed desorption pressure, simply depending on whether the relative change in Δq_CH4 or Δq_CO2 is greater as a function of desorption pressure (it was very rare for the relative changes to become equal in the desorption pressure range allowed). This can be unrelated to the adsorbent’s performance, e.g., for an adsorbent very selective for CO2 one may find that the sorbent selection parameter is maximized at a high desorption pressure because negligible changes in Δq_CH4 have a greater impact on the working capacity ratio than significant changes in Δq_CO2. Thus, the sorbent selection parameter cannot be used if one is interested in finding optimal desorption conditions.

We similarly compared the SPP to the alternative metric of the adsorbent performance indicator (API) given by Wiersum et al.18

\[
\text{API} = \frac{(a_{CO2, CH4, \text{abs}} - 1) \Delta q_{CO2}}{Q_{w, CO2}} \tag{7}
\]

However, the adsorbent performance indicator is highly correlated with the sorbent selection parameter, and suffers from much the same problems (Fig. S1†).

3. Results and discussion

3.1 Screening

We first investigated the relative merits of PSA and TSA by examining the optimal desorption pressure and desorption temperature for the IZA zeolites undergoing PTSA processes. For all three process, less than 10% of the zeolites displayed an optimal desorption temperature above 305 K, with most

Table 1 A selection of hypothetical zeolites with deviating SPP and sorbent selection parameter (S) values, shown in order of increasing Δq_CO2. The first three zeolites have smaller SPP but larger sorbent selection parameter than the last two zeolites. The values shown are for materials undergoing the same process conditions as in Fig. 1, where these zeolites are circled in green. Note that PCOD8238989 exhibits a higher SPP than PCOD8310046 despite having a higher Δq_CH4 and a lower Δq_CO2, which is due to it having a lower Q_w,CO2.

| Adsorbent         | Δq_CH4 (mol kg⁻¹) | Δq_CO2 (mol kg⁻¹) | SPP (mol² kg⁻¹ kJ⁻¹) | S     |
|-------------------|-------------------|-------------------|---------------------|-------|
| PCOD8294501      | 1 × 10⁻⁵          | 0.005             | 0.01                | 3.6 × 10⁷ |
| PCOD8164653      | 8 × 10⁻⁶          | 0.063             | 0.16                | 1.5 × 10⁷ |
| PCOD8306302      | 4 × 10⁻⁷          | 0.192             | 0.54                | 2.3 × 10⁷ |
| PCOD8238989      | 7 × 10⁻²          | 2.359             | 8.18                | 9.9 × 10⁴ |
| PCOD8310046      | 3 × 10⁻²          | 2.550             | 7.71                | 6.6 × 10⁴ |
materials benefiting more from pulling vacuum than from a temperature-swing (Fig. S2†). For the few zeolites that had a higher optimal desorption temperature, PTSA only negligibly increased the SPP from its value with PSA, so the temperature-swing does not appear to be truly beneficial for any material; conversely, PTSA significantly increased the SPP of several materials compared to that obtained with TSA, demonstrating that pulling vacuum can be truly beneficial (Fig. S3†). In industrial practice as well, PSA is generally preferable to TSA due to a shorter cycle time and lower complexity.6 It was thus concluded that only PSA processes need be analyzed for the larger datasets studied herein.

PSA screening was then performed for the three processes using all IZA and hypothetical zeolites. Optimal desorption pressures and the resulting SPP are given for the IZA zeolites in Tables S1–S3,† from which a list of top-performing materials can be identified. For the hypothetical zeolites, we plot the SPP as a function of $Q_{st,\text{CO}_2}$ and the CO$_2$ saturation loading capacity (as obtained by fitting the pure-component isotherm) in Fig. 2, and we illustrate topologies and potential energy surfaces of some of the highest-performing frameworks in Fig. 3. The LFG, APG, and NAG processes exhibit different ranges of optimal $Q_{st,\text{CO}_2}$ centered around approximately 37, 35 and 32 kJ mol$^{-1}$ respectively, decreasing with an increase in adsorption pressure as predicted by theory.34 Although these optimal $Q_{st,\text{CO}_2}$ values are different from each other, they are not so different that the zeolites’ SPP for different processes are not well-correlated with each other (Fig. S4†). However, this finding of zeolite performance being correlated between processes may be specific to our choice of processes, which exhibit some overlap in the CO$_2$ partial pressures over which they cycle; processes with a smaller overlap may in fact exhibit unique optimal adsorbents.8,11 Our optimal $Q_{st,\text{CO}_2}$ values are larger than those found in a screening of a large database of hypothetical metal–organic frameworks (MOFs) by Wilmer et al.,11 where processes similar to the LFG and APG processes showed optimal sorbent selection parameter with materials having $Q_{st,\text{CO}_2}$ values centered around approximately 29 and 31 kJ mol$^{-1}$, respectively. This is likely due in part to the different performance metric, pressures, and mole fractions used between our studies, and in part due to zeolites being more confining materials than MOFs, with typically smaller pores and void fractions35 that provide a larger decrease in entropy upon adsorption of CO$_2$ and thus a larger optimal $Q_{st,\text{CO}_2}$.34

Within the optimal ranges of $Q_{st,\text{CO}_2}$, the best materials are those with a sufficiently large CO$_2$ saturation loading capacity. The range of the SPP for all materials spans several orders of magnitude; by selecting only those materials with near-optimal $Q_{st,\text{CO}_2}$ values, the range is reduced but still spans orders of magnitude, and by filtering out the materials with too low a CO$_2$ saturation loading capacity the variability is reduced further. As others have found in large database screenings, bounds on the optimal performance of materials appear to follow simple functions of descriptors, but a given material’s performance cannot be perfectly predicted with these same descriptors.11,13

3.2 Performance prediction

One wishes to find material descriptors that allow one to predict a zeolite’s performance. For this purpose there exist
several well-established geometric descriptors of porous media including void fraction (as measured using a helium atom probe), crystal density, accessible surface area, largest included sphere diameter (largest sphere that can fit inside

![Fig. 3 Visualization of the (top) CH₄ and (bottom) CO₂ potential energy surfaces of some of the top-performing hypothetical zeolites for the APG PSA process. From left to right, the zeolites shown are PCOD8284104, PCOD8246225, PCOD8184913, PCOD8246205, and PCOD8310046. Contours correspond to potential energies of −30 kJ mol⁻¹ (purple), −10 kJ mol⁻¹ (blue), 0 kJ mol⁻¹ (green), and 37.4 kJ mol⁻¹ (gray), with the last representing the limits of accessible channels, approximately defined as 15R × 300 K, which our GPU GCMC code uses to determine blocked pockets inaccessible from the gas phase, which are then eliminated from use as possible adsorption sites. The second and fourth zeolites from the left demonstrate some CH₄ blocked pockets. Both cage- and channel-containing topologies are seen in these top-performing materials. Framework silicon and oxygen atoms are beige and red. The CO₂ potential energy shown at each grid point is a Boltzmann-weighted average of 1000 random rotations of the CO₂ molecule about the carbon atom which remains stationary at the grid point.]

![Molecular Systems Design & Engineering Paper](https://example.com)
the material), and largest free sphere diameter (largest sphere that can diffuse through the material). In addition, the pure-component isotherm descriptors $Q_{st}$ and saturation loading capacity are available for each adsorbate. To determine which of these descriptors are the most important predictors of an adsorbent's performance, we used a random forest machine learning algorithm, a method that Simon et al. used to screen similar materials for a Xe/Kr separation process and which provides a quantitative measure of a descriptor's importance to improving the quality of prediction. Since it has been shown that the helium void fraction and the crystal density are strongly negatively correlated while the largest included sphere diameter and the largest free sphere diameter are strongly positively correlated, we did not include the crystal density and the largest free diameter as descriptors in the regression; we verified that switching these choices did not give qualitatively different results.

Results from performing the random forest regression on all three processes are shown in Fig. 4. The CO$_2$ isotherm descriptors show the greatest influence. Fig. 2 and 4 both show that for the high-pressure NAG process, the SPP becomes less dependent on $Q_{st,CO_2}$ and more on the CO$_2$ saturation loading capacity, which can be easily rationalized based on the shape of a Langmuir isotherm, which is more sensitive to the Henry coefficient (and thus $Q_s$) at low pressures and saturation loading capacity at high pressures.

That the isotherm descriptors proved more important than the geometric descriptors is not altogether surprising. Separation performance is directly dictated by a material's adsorption isotherms, which in turn are determined by a zeolite's topology; thus, the geometric descriptors are at least one step further from separation performance than isotherm descriptors. Less intuitive is the seemingly minor role taken by the CH$_4$ isotherm descriptors. The low importance of the CH$_4$ isotherm descriptors cannot be attributed to the presence of a correlation with the CO$_2$ isotherm descriptors corrupting the regression results, as scatter plots show no clear relationship between CH$_4$ isotherm descriptors and the SPP, while the CO$_2$ isotherm descriptors are seen to create upper and lower limits on an adsorbent's performance (Fig. S7–S9†).

The effect of the CH$_4$ isotherm descriptors' influence on separation performance is complicated by the fact that the CH$_4$ and CO$_2$ isotherm descriptors are correlated (Fig. 5). While saturation loading capacities of the two species may be difficult to decouple due their heavy dependence on an adsorbent's topology, $Q_{st,CH_4}$ and $Q_{st,CO_2}$ can be separately tuned by taking advantage of the molecules' difference in electrostatics. It would therefore be useful to observe the effect of removing this correlation and then varying the $Q_{st}$ of one adsorbate while keeping the $Q_{st}$ of the other adsorbate constant. We accomplished this by varying the guest–host force field parameters for one adsorbate at a time, i.e. to see the individual effect of $Q_{st,CH_4}$, the CH$_4$–O$_{zeo}$ epsilon parameter of the Lennard-Jones potential was varied above and below its original value, while the individual effect of $Q_{st,CO_2}$ was observed by varying the C$_{CO_2}$–O$_{zeo}$ and O$_{CO_2}$–O$_{zeo}$ epsilon parameters. The resulting $Q_{st}$ are shown along with the SPP of the PSA APG process for a subset of the IZA zeolites in Fig. 6. It can be seen that lowering $Q_{st,CO_2}$ independently of $Q_{st,CO_2}$ increases performance for a range of zeolites nearly uniformly, with the SPP approaching limiting values where no CH$_4$ adsorbs at all. Conversely, $Q_{st,CO_2}$ exhibits an optimal value unique to each zeolite and process, with performance getting worse as $Q_{st,CO_2}$ decreases or increases from this optimal value. For several of the zeolites, the $Q_{st,CO_2}$ graph is broadly peaked, with the SPP remaining near its maximum value over a range of about 10 kJ mol$^{-1}$. The top-performing zeolites seem to exhibit a $Q_{st,CO_2}$ near their optimal values, but there clearly exist materials which would perform poorly even at their optimal $Q_{st,CH_4}$ due to a low CO$_2$ saturation loading capacity. Overall, an ideal material with uncorrelated heats of adsorption for CH$_4$ and CO$_2$ would exhibit a very small $Q_{st,CH_4}$ and an appropriately-sized $Q_{st,CO_2}$.

To understand these trends, consider the argument of Bhatia and Myers that the working capacity of a single-component gas in a porous material is maximized at a particular $Q_{st}$, which is simply demonstrated by finding the analytical maximum of the difference between loadings at two pressures of a Langmuir isotherm, and positing that the entropy of adsorption is similar in many materials. One expects a similar argument to hold for CO$_2$ in our multicomponent separation, and one might also expect that to maximize the

![Fig. 4 Importance of geometric and isotherm descriptors to the SPP of the hypothetical zeolites undergoing PSA processes, as determined using a random forest of decision trees. The importance of a descriptor is calculated by summing the reductions in mean squared error brought about at each node where that descriptor splits a decision tree, averaging over all decision trees, and normalizing. Here, sat. load. is the saturation loading capacity, $Q_s$ is the largest included sphere diameter, Accessible SA is the accessible surface area, and He void frac. is the helium void fraction.](image-url)
SPP, one would like a $Q_{\text{st,CH}_4}$ that is either much larger or much smaller than the $Q_{\text{st,CH}_4}$ that maximizes its single-component working capacity. Since a large $Q_{\text{st,CH}_4}$ takes adsorption sites away from CO$_2$, it instead makes sense that the SPP increases with a decreasing $Q_{\text{st,CH}_4}$. Indeed, even for the zeolite FER, in which SPP is seen to increase with increasing $Q_{\text{st,CH}_4}$ beyond a certain point, the CO$_2$ working capacity is decreasing, and the improved SPP is due only to the working
capacity of CH₄ decreasing at an even faster rate than that of CO₂. Notably, the zeolites do not exhibit a single optimal \( Q_{st,CO_2} \) but rather a range of optimal \( Q_{st,CO_2} \) values, e.g., for the APG process, the zeolites STI, WEI, and MOR exhibit optimal \( Q_{st,CO_2} \) of 32, 41, and >50 kJ mol⁻¹, respectively, reflecting the negation of the assumption of similar entropies of adsorption across all materials, which Simon et al.⁴² has recently demonstrated.

Since \( Q_{st,CH_4} \) and \( Q_{st,CO_2} \) are positively correlated, it becomes difficult to obtain an optimal \( Q_{st,CO_2} \) while minimizing \( Q_{st,CH_4} \), so \( Q_{st,CH_4} \) becomes difficult to relate to a zeolite’s performance, which explains the low importance of \( Q_{st,CH_4} \), obtained from the regression. This understanding suggests that the ability to tune the CO₂-host interaction potential independently of the CH₄-host interaction potential can lead to higher-performing materials; this can be achieved in practice by changing the charges of framework atoms, more easily done with different material classes like cation-exchanged zeolites or MOFs.

3.3 Structure–property relationships

To aid in the design of new materials, it is preferable to be able to predict a material’s performance based solely on geometric descriptors (e.g., helium void fraction, crystal density, accessible surface area, largest included sphere diameter, and largest free sphere diameter), and so we would like to eliminate the use of isotherm descriptors. Within a material class, one would expect this to be possible in theory, particularly for purely-siliceous zeolites as they have identical stoichiometric makeup. Since we have shown that isotherm descriptors are well-capable of predicting a zeolite’s performance, if geometric proxies can be found that have a strong relationship with the important isotherm descriptors, the necessity for the isotherm descriptors should be able to be eliminated. However, while it is known that saturation loading capacities of both CO₂ and CH₄ can be predicted by several simple geometric descriptors, and past studies have shown a clear relationship between a material’s largest sphere diameter and CH₄ adsorption energetics,⁴³ more complex descriptors are necessary to predict \( Q_{st,CO_2} \) (Fig. S13 and S14†). Since we lack a simple geometric proxy for \( Q_{st,CO_2} \), we will be unable to remove it from use. Indeed, although past screening studies have shown that purely geometric descriptors can predict performance for single-component gas storage of simple molecules like CH₄,⁴⁶ for more complicated multicomponent gas separation processes involving CO₂ it has been found necessary for geometric parameters to be supplemented by \( Q_{st,CO_2} \) or by both \( Q_{st} \) values.¹¹,¹⁷ It appears that the inclusion of \( Q_{st,CO_2} \) (or a similar isotherm descriptor such as the Henry coefficient) as a performance predictor will be necessary until simple geometric descriptors that are well-correlated with \( Q_{st,CH_4} \) are found.

Conversely, one would expect to find the CO₂ saturation loading capacity unnecessary because it seems to have strong relationships with several geometric descriptors. However, the relationship between the geometric descriptors and the SPP is more ambiguous than that between the CO₂ saturation loading capacity and the SPP, so a simple heuristic cannot be deduced (Fig. S15–S20†), though the descriptor is less crucial for machine learning prediction.

3.4 IAST validity

As we have generated both pure and mixture adsorption isotherms for over 80 000 zeolites, we can study the validity of the IAST assumption for nanoporous materials at a larger scale than has been attempted before. Many studies have evaluated the validity of IAST for a small number of materials.⁴⁸–⁵⁷ Of these, the ones that have focused on mixtures of CO₂ and CH₄ in zeolites and MOFs have generally found good agreement between IAST and mixture isotherms at low pressures but some divergence emerging at pressures above about 1 bar, with the various studies showing that divergence can manifest itself in several ways.⁴⁸–⁵²

One reason commonly cited for the violation of IAST is its assumption that the same adsorbent surface area is available to all adsorbate molecules,⁵³,⁵⁴ an assumption that Myers and Prausnitz²⁷ acknowledged is invalid for molecular sieves, and which tends to be less valid at higher pressures as adsorbate molecules compete for adsorption sites and adsorbates become accommodated at less energetically favorable adsorption sites. Another reason commonly cited for the violation of IAST is its assumption of adsorbed-phase ideality,⁵⁵ which also becomes less valid at higher pressures as higher adsorbate loadings are induced.

We compared mixture isotherm loadings obtained directly with those obtained through IAST, with results shown for a 60 mol% CH₄ mixture in Fig. 7 and qualitatively similar results shown for a 90 mol% CH₄ mixture in Fig. S21†. For CH₄, it can be seen that at low pressures the IAST error is centered around 0, but as the pressure increases IAST tends to underpredict the loading, with errors becoming considerable by around 5 bar. For CO₂, IAST brings about less bias at high pressures as adsorbate molecules compete for adsorption sites and adsorbates become accommodated at less energetically favorable adsorption sites. These divergence trends are similar to what was seen in Krishna and van Baten,⁵⁰ where it was found that different adsorption sites (e.g., windows versus cages) contain different proportions of the two adsorbates; this leads IAST to anticipate greater competition between the adsorbates than actually occurs, and the loading of the weaker-adsorbing species is underpredicted. This type of error brought about by IAST due to site segregation can be eliminated by using a segregated IAST model.⁵⁴

For both adsorbates and at all pressures, there exist materials in which IAST introduces large errors. One of the materials with the largest introduced error, PCOD8205017, was investigated in further detail. Pure-component and mixture isotherms of CH₄ and CO₂ in this material are shown in Fig. 8, where it can be seen that the directly-generated mixture isotherms are similar to the pure-component isotherms, and that IAST drastically underpredicts CH₄ loading. The snapshots in Fig. 8 show that the CH₄ and CO₂ adsorbates

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occupy very distinct adsorption sites; these same sites are also preferentially occupied by each adsorbate during pure-component adsorption, so during mixture adsorption there is a lack of competition between adsorbates for adsorption sites until a pressure is reached at which one of the two sites becomes saturated. This extreme case of site segregation illustrates the dangers of not verifying the validity of IAST, as IAST results lead to the prediction that PCOD8205017 is a much higher-performing material than it actually is.

By further examining the PCOD8205017 isotherms in Fig. 8, it can be seen that at certain points the uptakes are even larger in the directly-generated mixture isotherms than in the pure-component isotherms. This is due to the CH₄–CO₂ and CH₄–OCO₂ adsorption site distances both being approximately 1.1 Å.

Fig. 7 Loadings of (a) CH₄ and (b) CO₂ at various total pressures, 300 K, and 60 mol% CH₄. Data on the x-axis are taken from directly-simulated mixture isotherms, while data on the y-axis are taken from applying IAST to pure-component isotherms. A line is drawn at y = x for reference. Hypothetical zeolites are shown as black dots, IZA zeolites are shown as blue diamonds, and hypothetical zeolite PCOD8205017 is shown as a green circle. Histograms of the fractional IAST error of the hypothetical zeolites for (c) CH₄ and (d) CO₂ at the same conditions. A line is drawn at x = 0 for reference. All plots in the left-column are at 0.1 bar, all plots in the middle-column are at 1 bar, and all plots in the right-column are at 5 bar. A similar plot showing results at 90 mol% CH₄ is given in Fig. S21.
times their Lennard-Jones sigma values, near the potential’s minimum. Thus, rather than following the IAST assumption that adsorbates compete for the same adsorption sites, this guest–guest nonideality instead makes the presence of the second adsorbate species conducive to the uptake of the first. That only small deviations from ideality occur despite such positioning of adsorption sites lends support to the ideality assumption of IAST segregated IAST, as well as to a previous finding that guest–host interactions give rise to nonideality to a greater extent than do guest–guest interactions.\textsuperscript{58}

For the purposes of high-throughput materials screening, a small number of outliers is less important than the introduction of systematic errors. To test whether IAST brought about systematic errors, we reid our earlier analysis of descriptors using the IAST results, and we obtained qualitatively-similar results, with the higher-pressure NAG process associated with a decrease in the optimal $Q_{st,CO_2}$ values and an increased dependence on the CO$_2$ saturation loading capacity (Fig. S22 and S23\textsuperscript{†}). Thus, IAST appears to be adequate for the purposes of performing high-throughput materials screening to find trends that relate structure to performance, but it must be tested for validity when performing detailed analyses of individual adsorbents.

4. Conclusions

In this work, we have developed the SPP metric, which avoids the pitfalls associated with other adsorption metrics such as the sorbent selection parameter and the adsorbent performance indicator: namely that they can overemphasize negligible gains in selectivity at the expense of more important changes in working capacity. Although the SPP was developed specifically for a process in which CO$_2$ is adsorbed from a CH$_4$/CO$_2$ mixture, it can be easily adapted for other adsorption processes.

The screening work that we then undertook with the SPP is complementary to several recent screening analyses of materials for CO$_2$ capture.\textsuperscript{8,11–13,47–61} Similarly to several of these studies\textsuperscript{8,11,13} we found that $Q_{st,CO_2}$ should be neither too low nor too high to make an optimal material. We built on this finding to show that the CO$_2$ saturation loading capacity is another important parameter to consider when designing materials, particularly for high-pressure processes. We then extended the theory of Bhatia and Myers,\textsuperscript{34} initially developed for storage of a single-component gas, to a multicomponent gas separation case by demonstrating that $Q_{st,CH_4}$ should be as low as possible to optimize separations, but that because $Q_{st,CH_4}$ is positively correlated with $Q_{st,CO_2}$, it misleadingly appears to be an unimportant parameter.

Since we chose to directly generate mixture isotherms rather than rely on the predictions of IAST, we were also able to perform a large-scale evaluation of the validity of IAST for siliceous zeolites. We found that IAST works reliably for the purposes of using high-throughput screening to find performance trends, but above about 5 bar IAST begins to underpredict the loading of CH$_4$. At all pressures there exist zeolites for which IAST does not work, which we showed can be due to site segregation.

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References

1 Annual Energy Outlook 2015, U. S. Energy Information Administration technical report, 2015.
2 S. Pacala and R. Socolow, Science, 2004, 305, 968–972.
3 R. W. Howarth, R. Santoro and A. Ingraffea, Clim. Change, 2011, 106, 679–690.
4 A. R. Brandt, G. A. Heath, E. A. Kort, F. O’Sullivan, G. Pétron, S. M. Jordaan, P. Tans, J. Wilcox, A. M. Gopstein, D. Arent, S. Woofy, N. J. Brown, R. Bradley, G. D. Stucky, D. Eardley and R. Harriss, Science, 2014, 343, 733–735.
5 N. Kumar, P. Besuner, S. Lefton, D. Agan and D. Hilleman, Power Plant Cycling Costs, NREL subcontract report prepared by Intertek API TECH, 2012.
6 M. Tagliabue, D. Farrusseng, S. Valencia, S. Aguado, U. Ravon, C. Rizzo, A. Corma and C. Mirodatos, Chem. Eng. J., 2009, 155, 533–566.
7 S. Cavenati, C. A. Grande and A. E. Rodrigues, J. Chem. Eng. Data, 2004, 49, 1095–1101.
8 J. M. Simmons, H. Wu, W. Zhou and T. Yildirim, Energy Environ. Sci., 2011, 4, 2177–2185.
9 Y. Bae and R. Q. Snurr, Angew. Chem., Int. Ed., 2011, 50, 11586–11596.
10 R. Krishna and J. M. van Baten, Phys. Chem. Chem. Phys., 2011, 13, 10593–10616.
11 C. E. Wilmer, O. K. Farha, Y. Bae, J. T. Hupp and R. Q. Snurr, Energy Environ. Sci., 2012, 5, 9849–9856.
12 J. Kim, A. Maiti, L.-C. Lin, J. K. Stolaroff, B. Smit and R. D. Aines, Nat. Commun., 2013, 4, 1694.
13 L.-C. Lin, A. H. Berger, R. L. Martin, J. Kim, J. A. Swisher, K. Jariwala, C. H. Rycroft, A. S. Bhown, M. W. Deem, M. Haranczyk and B. Smit, Nat. Mater., 2012, 11, 633–641.
14 B. J. Maring and P. A. Webley, Int. J. Greenhouse Gas Control, 2013, 15, 16–31.
15 M. M. F. Hasan, E. L. First and C. A. Floudas, Phys. Chem. Chem. Phys., 2013, 15, 17601–17618.
16 E. L. First, M. M. F. Hasan and C. A. Floudas, AIChE J., 2014, 60, 1767–1785.
17 S. Rege and R. Yang, Sep. Sci. Technol., 2001, 36, 3355–3365.
18 A. D. Wiersum, J.-S. Chang, C. Serre and P. L. Llewellyn, Langmuir, 2013, 29, 3301–3309.
19 C. Baerlocher and L. B. McCusker, Database of Zeolite Structures, http://www.iza-structure.org/databases/, Online; accessed Jan. 1, 2010.
20 R. Pophale, P. A. Cheeseman and M. W. Deem, Phys. Chem. Chem. Phys., 2011, 13, 12407.
21 J. Kim, R. L. Martin, O. Rübel, M. Haranczyk and B. Smit, J. Chem. Theory Comput., 2012, 8, 1684–1693.
22 J. Kim and B. Smit, J. Chem. Theory Comput., 2012, 8, 2336–2343.
23 D. Frenkel and B. Smit, Understanding Molecular Simulation: From Algorithms to Applications, Elsevier Science, 2002.
24 E. García-Pérez, J. B. Parra, C. O. Ania, A. García-Sánchez, J. M. van Baten, R. Krishna, D. Dubbeldam and S. Calero, Adsorption, 2007, 13, 469–476.
25 D. Peng and D. B. Robinson, Ind. Eng. Chem. Fundam., 1976, 15, 59–64.
26 O. Talu and A. L. Myers, AIChE J., 2001, 47, 1160–1168.
27 A. L. Myers and J. M. Prausnitz, AIChE J., 1965, 11, 121–127.
28 E. Jones, T. Oliphant and P. Peterson, et al., Scipy: Open source scientific tools for Python, 2001, http://www.scipy.org/, Online; accessed Apr. 8, 2015.
29 H. Akima, J. Assoc. Comput. Mach., 1970, 17, 589–602.
30 J. Boerio-Goates, R. Stevens, B. K. Hom, B. F. Woodfield, P. M. Piccione, M. E. Davis and A. Navrotsky, J. Chem. Thermodyn., 2002, 34, 205–227.
31 S. Krishnamurthy, V. R. Rao, S. Guntaka, P. Sharratt, R. Haghpanah, A. Rajendran, M. Amanullah, I. A. Karimi and S. Farooq, AIChE J., 2014, 60, 1830–1842.
32 J. M. Smith, H. C. Van Ness and M. M. Abbott, Introduction to Chemical Engineering Thermodynamics, McGraw-Hill, 7th edn, 2005.
33 E. Braun, A. F. Zurhelle, W. Thijssen, S. K. Schnell, L.-C. Lin, J. Kim, J. A. Thompson and B. Smit, High-Throughput Computational Screening of Nanoporous Adsorbents for CO2 Capture from Natural Gas, Open Science Framework, 2016, DOI: 10.17605/OSF.IO/KTBV4, Online.
34 S. K. Bhatia and A. L. Myers, Langmuir, 2006, 22, 1688–1700.
35 C. M. Simon, J. Kim, D. A. Gomez-Gualdron, J. S. Camp, Y. G. Chung, R. L. Martin, R. Mercado, M. W. Deem, D. Gunter, M. Haranczyk, D. S. Sholl, R. Q. Snurr and B. Smit, Energy Environ. Sci., 2015, 8, 1190–1199.
36 T. F. Willems, C. H. Rycroft, M. Kazi, J. C. Meza and M. Haranczyk, Microporous Mesoporous Mater., 2012, 149, 134–141.
37 M. Pinheiro, R. L. Martin, C. H. Rycroft and M. Haranczyk, CrystEngComm, 2013, 15, 7531–7538.
38 L. Breiman, Mach. Learn., 2001, 45, 5–32.
39 F. Pedregosa, G. Varoquaux, A. Gramfort, V. Michel, B. Thirion, O. Grisel, M. Blondel, P. Prettenhofer, R. Weiss, V. Dubourg, J. Vanderplas, A. Passos, D. Cournapeau, M. Brucher, M. Perrot and E. Duchesnay, J. Mach. Learn. Res., 2011, 12, 2825–2830.
40 C. M. Simon, R. Mercado, S. K. Schnell, B. Smit and M. Haranczyk, Chem. Mater., 2015, 27, 4459–4475.
41 L. Breiman, J. Friedman, C. J. Stone and R. Olshen, Classification and Regression Trees, Chapman & Hall/CRC, 1984.
42 C. M. Simon, J. Kim, L.-C. Lin, R. L. Martin, M. Haranczyk and B. Smit, Phys. Chem. Chem. Phys., 2014, 16, 5499–5513.
43 E. Haldoupis, S. Nair and D. S. Sholl, *Phys. Chem. Chem. Phys.*, 2011, 13, 5053–5060.
44 H. Amrouche, B. Creton, F. Siperstein and C. Nieto-Draghi, *RSC Adv.*, 2012, 2, 6028–6035.
45 E. J. Garcia, J. Pérez-Pellitero, C. Jallut and G. D. Pirngruber, *Phys. Chem. Chem. Phys.*, 2013, 15, 5648–5657.
46 M. Fernandez, T. K. Woo, C. E. Wilmer and R. Q. Snurr, *J. Phys. Chem. C*, 2013, 117, 7681–7689.
47 D. Wu, Q. Yang, C. Zhong, D. Liu, H. Huang, W. Zhang and G. Maurin, *Langmuir*, 2012, 28, 12094–12099.
48 R. Babarao, Z. Hu, J. Jiang, S. Chempath and S. I. Sandler, *Langmuir*, 2007, 23, 659–666.
49 H. Chen and D. S. Sholl, *Langmuir*, 2007, 23, 6431–6437.
50 R. Krishna and J. M. van Baten, *Sep. Purif. Technol.*, 2008, 61, 414–423.
51 Y.-S. Bae, K. L. Mulfort, H. Frost, P. Ryan, S. Punnathanam, L. J. Broadbelt, J. T. Hupp and R. Q. Snurr, *Langmuir*, 2008, 24, 8592–8598.
52 S. E. Jee and D. S. Sholl, *J. Am. Chem. Soc.*, 2009, 131, 7896–7904.
53 N. F. Cessford, N. A. Seaton and T. Düren, *Ind. Eng. Chem. Res.*, 2012, 51, 4911–4921.
54 J. A. Swisher, L.-C. Lin, J. Kim and B. Smit, *AIChE J.*, 2013, 59, 3054–3064.
55 A. Goj, D. S. Sholl, E. D. Akten and D. Kohen, *J. Phys. Chem. B*, 2002, 106, 8367–8375.
56 E. D. Akten, R. Siriwardane and D. S. Sholl, *Energy Fuels*, 2003, 17, 977–983.
57 M. Murthi and R. Q. Snurr, *Langmuir*, 2004, 20, 2489–2497.
58 M. Sakuth, J. Meyer and J. Gmehling, *Chem. Eng. Process.: Process Intensif.*, 1998, 37, 267–277.
59 E. Haldoupis, S. Nair and D. S. Sholl, *J. Am. Chem. Soc.*, 2012, 134, 4313–4323.
60 J. Kim, L.-C. Lin, J. A. Swisher, M. Haranczyk and B. Smit, *J. Am. Chem. Soc.*, 2012, 134, 18940–18943.
61 J. Kim, M. Abouelnasr, L.-C. Lin and B. Smit, *J. Am. Chem. Soc.*, 2013, 135, 7545–7552.