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Physical Properties of XN (X = B, Al, Ga, In) in the Pm–3n phase: First-Principles Calculations

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Abstract: Three direct semiconductor materials and one indirect semiconductor material, Pm–3n XN (X = B, Al, Ga, In), are investigated in our work, employing density functional theory (DFT), where the structural properties, stability, elastic properties, elastic anisotropy properties and electronic properties are included. The shear modulus $G$ and bulk modulus $B$ of Pm–3n BN are 290 GPa and 244 GPa, respectively, which are slightly less than the values of $B$ and $G$ for c-BN and Pnma BN, while they are larger than those of C$_6$4 in the I4_1/amd phase. The shear modulus of Pm–3n BN is the greatest, and the shear modulus of C$_6$4 in the I4_1/amd phase is the smallest. The Debye temperatures of BN, AlN, GaN and InN are 1571, 793, 515 and 242 K, respectively, using the elastic modulus formula. AlN has the largest anisotropy in the Young’s modulus, shear modulus, and Poisson’s ratio; BN has the smallest elastic anisotropy in $G$; and InN has the smallest elastic anisotropy in the Poisson’s ratio. Pm–3n BN, AlN, GaN and InN have the smallest elastic anisotropy along the (111) direction, and the elastic anisotropy of the $E$ in the (100) (010) (001) planes and in the (011) (101) (110) planes is the same. The shear modulus and Poisson’s ratio of BN, AlN, GaN and InN in the Pm–3n phase in the (001), (010), (100), (111), (101), (110), and (011) planes are the same. In addition, AlN, GaN and InN all have direct band-gaps and can be used as a semiconductor within the HSE06 hybrid functional.

Keywords: III-V nitride compounds; elastic anisotropy; direct band gap; stability

1. Introduction

In the 1950s, Germanium was used in low-voltage, low-frequency, medium power transistors and photodetectors in the bright stage. However, there were large short plates in the high-temperature and -radiation resistance of germanium semiconductor devices, so, in the 1960s, germanium gave up its dominant position to silicon. Silicon diffusely serves as a semiconductor material in industry and is mainly used in data computing and other fields. With the increasing demands of science and technology, the defects in the slow transmission speed and single function of silicon have been exposed, so compound semiconductor materials have emerged as required. Third-generation semiconductor materials have many ascendancies, such as a broad band gap, large thermal conductivity, great electron saturation rate, huge breakdown electric field, and a strong ability to resist radiation, so they have a wide range of applications in some blue, green, and violet light-emitting diodes and semiconductor lasers. GaN has a faster speed and higher breakdown voltage compared to silicon. Compared with silicon devices, GaN is more advanced in power conversion efficiency and power density. AlN is an ideal substrate material for advanced high-power light-emitting devices (LEDs, LDs), UV probes and high-power high-frequency electronic equipments. Si and compound semiconductors are two complementary materials. Some of the performance advantages of the compounds make up
for the shortcomings of Si crystals, while the production process of Si crystals has obvious irreplaceable advantages, and both of them have certain limitations in the field of application. Therefore, in the application of semiconductors, compatibility means are often used to make the two compatible and utilize their respective advantages. Because of this, more and more attention has been paid to research on new structures and the physical properties of the compounds [1–13].

Lately, a growing number of studies have been conducted in the field of novel semiconductor materials, such as III-V nitride compound [1–9], other III-V compound materials [13–20], carbon-based [14–21], and silicon-based [22–28]. The structural properties, electronic properties, mechanical attributes, and stableness of the BN polymorph in the Pnma structure were investigated, utilizing first-principles calculations by the Cambridge Serial Total Energy Package (CASTEP) plane-wave code, which was studied by Ma et al [1]. They also discovered that Pnma-BN has larger band gap (7.18 eV) than other BN polymorphs, and it has an indirect band gap with the HSE06 function [29]. Fan et al. [3] reported the physical properties of AlN, GaN, and InN in the Pnma phase. Different from Pnma-BN, all the AlN, GaN, and InN materials in the Pnma phase have direct band gaps, and Pnma-AlN, Pnma-GaN, and Pnma-InN are the latent binary semiconductor materials for the production of UV detectors, violet photodiodes and infrared detectors, respectively. Compared with the previous materials (AlGaAs, GaAsP, AlGaN, AlGaN and other ternary semiconductor alloys, or AlGaInP quaternary semiconductor alloy), they do not require doping [3]. Liu et al [5] proposed four new AlN polymorphs, named Pbca-AlN, Pbca-AIN, Pmnn2-AlN and Cmcm-AlN, and Yang et al [6,7] studied the physical properties of six new AlN polymorphs, Pnn21-AlN, Cmcm-AlN, Pbca-AlN, Pbam-AlN, bct-AlN (space group: P42/mnm), and h-AlN (space group: P63/mmc). All the AlN polymorphs, Pnn21-AlN, Pbam-AlN, Pbca-AlN, Cmcm-AlN, P42/mnm-AlN, are direct band gap semiconductor materials, and their band gaps are 3.63 (Cmcm-AlN), 3.89 (Pnn21-AlN), 3.86 (Pbca-AlN), 3.93 (Pbam-AlN), and 5.85 eV (P42/mnm-AlN), with Perdew–Burke–Ernzerhof (PBE) functionality and PBE0 functionality, respectively.

In view of the physical properties of the whole III-V nitrides of adamantine phase studied by density functional theory [1–3], the band structure of AlN, Gan and InN is direct band gap, which has great potential application in the photoelectric industry or dye-sensitive solar cells [3]. So, according to density functional theory [30,31], the III-V nitride compounds, BN, AlN, GaN, InN in the Pm–3n phase are first proposed in our paper. The physical attributes of Pm–3n XN (X = B, Al, Ga, In) semiconductor materials are estimated and discussed, and the stability of Pm–3n XN (X = B, Al, Ga, In) is also investigated in this work. Additionally, their dynamic and mechanical stability are proven by phonon spectra, elastic constants and related enthalpies.

2. Theoretical Methods

Physical property predictions and structural geometric optimization calculations utilize DFT with ultrasoft pseudopotentials [32] under the Cambridge Sequential Total Energy Package (CASTEP) [33] code in Materials Studio. The exchange correlation potentials are used with the Perdew–Burke–Ernzerhof (PBE) functional of the generalized gradient approximation (GGA) [34], and the Broyden–Fletcher–Goldfarb–Shanno (BFGS) [35] minimization scheme adopts the geometric optimization of the Pm–3n XN (X = B, Al, Ga, In). A fairly high k-point separation (less than or ~ 0.025 Å⁻¹ × 2π) is applied to Pm–3n XN, 10 × 10 × 10, 8 × 8 × 8, 8 × 8 × 8, and 6 × 6 × 6 of the conventional cell for Pm–3n XN. Additionally, the plane-wave cutoff energy of 420 eV is adopted for structural optimizations and property predictions for Pm–3n XN (X = B, Al, Ga, In). The Heyd–Scuseria–Ernzerhof (HSE06) hybrid functional [29] was applied to the estimation of the electronic band structures. Finally, all the phonon spectra of Pm–3n XN adopt the density functional perturbation theory (DFPT) approach [36].
3. Results and Discussion

3.1. Structural Properties

The crystal textures of $Pm\overline{3}n$ $X_N$ ($X = B, Al, Ga, In$) are illustrated in Figure 1a. Here, red spheres represent nitrogen atoms, blue spheres represent group IIIB elements, and we mainly refer to boron atoms, aluminium atoms, gallium atoms and indium atoms in this work. Figure 1a shows that the quaternary ring structure composed of two nitrogen atoms and two atoms of the group IIIB elements is connected with four vertical quaternary rings. The quaternary ring structure is a square with nitrogen and boron (or aluminium, gallium or indium) atoms at the top of the square. In the conventional cell of $Pm\overline{3}n$ $X_N$, there is a 24-atom cage structure composed of nitrogen atoms and atoms of the group IIIB elements. As both c-BN and $Pm\overline{3}n$ InN belong to the cubic crystal system, the crystal structure of $Pm\overline{3}n$ InN has high symmetry, so the stacking mode of the atoms along the (010) and (100) sides is exactly the same. The related results are plotted in Figure 1b,c. The lattice constants of $Pm\overline{3}n$ InN obtained in this work are also credible.

![Crystal Structure](image)

**Figure 1.** The crystalline structures of $Pm\overline{3}n$ $X_N$ ($X = B, Al, Ga, In$) (a), along the (010) direction (b) and the (100) direction (c).

### Table 1. Lattice constants (Å), volumes of the conventional cell ($Å^3$), elastic parameters (GPa), $B$, $G$, $E$ (GPa) and Poisson’s ratio $v$ of $Pm\overline{3}n$ $X_N$ and c-BN.

| Materials | $a$  | $V$  | $C_{11}$ | $C_{12}$ | $C_{44}$ | $B$  | $G$  | $B/G$ | $E$  | $v$  |
|-----------|------|------|----------|----------|----------|------|------|-------|------|------|
| BN        | 4.438 | 87.416 | 700      | 85       | 209      | 290  | 244  | 1.189 | 572  | 0.171 |
| AlN       | 5.366 | 154.505 | 335      | 59       | 58       | 151  | 83   | 1.819 | 210  | 0.268 |
| GaN       | 5.584 | 174.088 | 238      | 61       | 58       | 120  | 69   | 1.739 | 174  | 0.259 |
| InN       | 6.237 | 242.570 | 173      | 55       | 36       | 95   | 44   | 2.159 | 114  | 0.299 |
| c-BN      | 3.622 | 47.517 | 779      | 165      | 446      | 370  | 384  | 0.964 | 856  | 0.115 |
|           | 3.620 $^a$ | 820 $^b$ | 190      | 480      | 400      |      |      |       |      |      |

$^a$ Reference [37]; $^b$ Reference [38].

3.2. Stability

Stability plays an extremely significant role in physical performance. The phonon spectra of $Pm\overline{3}n$ $X_N$ are plotted in Figure 2a–d. There is no frequency under the zero line, which proves that $Pm\overline{3}n$ $X_N$ is dynamically stable. In addition, its mechanical stability is also studied. The highest calculated phonon frequency of the B-N bond-stretching schema in $Pm\overline{3}n$ BN is ~38 THz, which is very close to that of diamond (40 THz), showing that the B-N bond in $Pm\overline{3}n$ BN is relatively strong. The elastic parameters are estimated through the strain stress method, as shown in Table 1. The three necessary and sufficient Born stability criteria for cubic symmetry are taken as: $C_{11} - C_{12} > 0$, $C_{11} + 2C_{12} > 0$, and $C_{44} > 0$. From Table 1, all the values for $C_{11}$, $C_{12}$ and $C_{44}$ of $Pm\overline{3}n$ $X_N$ satisfy the
three stability standards for cubic symmetry, which proves that \( Pm-3n \) \( XN \) is mechanically stable. Finally, the related enthalpies of \( Pm-3n \) \( XN \) are presented in Figure 2e–g. The enthalpy of \( XN \) of the wurtzite structure is set to 0. Although the enthalpy of \( Pm-3n \) BN is larger than that of \( Pnma \) BN [16], it is lower than that of rocksalt-BN and NiAs-BN [19]. For AlN, the enthalpy of the \( Pm-3n \) phase (0.177 eV/molecule) is slightly greater than that of \( Pbca \) AlN (0.172 eV/molecule) [20], while it is smaller than that of \( Pnma \) AlN (0.231 eV/molecule) [16]. Similar to \( Pm-3n \) AlN, the enthalpies of \( Pm-3n \) GaN (0.263 eV/molecule) and \( Pm-3n \) InN (0.226 eV/molecule) are also slightly smaller than those of the \( Pnma \) phase (\( Pnma \) GaN: 0.271 eV/molecule; \( Pnma \) InN: 0.237 eV/molecule).

Figure 2. Cont.
Although the value of \( B \) and ductility of Pm are the most brittle, whereas a lower \( B \) describes the plastic deformation opposition of the material, and \( E \) can be used to describe the tensile elasticity. Therefore, the ratio of \( B \) and \( G \) is a significant indicator to evaluate the brittleness or ductility of crystals. In accordance with Pugh [42], a higher \( B/G \) of a solid (more than 1.75) indicates improved ductility, whereas a lower \( B/G \) (less than 1.75) generally indicates brittleness. In addition, the value of \( v \) is in line with the \( B/G \), which touches on the ductility of the material; usually, the \( v \) value is large (\( v > 0.26 \)) [43]. The values of \( B/G \) and \( v \) for \( Pm-3n \) XN are plotted in Figure 3c. From Figure 3c, \( Pm-3n \) BN is the most brittle, and \( Pm-3n \) InN is the most ductile. \( Pm-3n \) GaN is in between the brittleness and ductility of \( Pm-3n \) BN and \( Pm-3n \) InN, but tends to be brittle.

Figure 3. The elastic parameters (a) and \( B, G \) and \( E \) (b) of \( Pm-3n \) XN and \( B/G \) and \( v \) of \( Pm-3n \) XN (c).
Knowledge of $\Theta_D$ is a powerful tool to reflect the bonding force between atoms. The value of $\Theta_D$ of different materials is different, and a large melting point indicates that the bond strength of the material is strong, then the Debye temperature increases as the Young’s modulus increases. The Debye temperature can be estimated by the elastic moduli formula. The formula are expressed as $v_p = [(B + 4G/3)/\rho]^{1/2}$, $v_s = (G/\rho)^{1/2}$, and $v_m = [(2/3\rho) + 1/v_p^2]/3^{1/3}$, where $\Theta_D = (h/k_B)[3n/(4\pi N_A \rho / M)]^{1/3} v_m$, $B$ and $G$ typifies the bulk modulus and shear modulus of AlN, GaN and InN, $n$ typifies the number of atoms in the molecule, $M$ typifies the molecular weight, $\rho$ typifies the crystal density, $h$ typifies Planck’s constant, $k_B$ represents Boltzmann’s constant, and $N_A$ represents Avogadro’s number. The crystal density, $v_s$, $v_p$ and $v_m$ and Debye temperature of AlN in the $Pm$–$3n$, $Pnma$ and $F$–$43m$ phases are shown in Table 2. For BN in the $Pm$–$3n$ and $Pnma$ phases, the Debye temperatures are close, while that of BN in the $Pm$–$3n$ phase is a little bit larger than that of BN in the $Pnma$ phase. Comparing the Young’s modulus, as shown in Table 1, $Pm$–$3n$ BN is a little bit higher than that of BN in $Pnma$ BN (543 GPa) [16]. For AlN, GaN and InN in the $Pm$–$3n$, $Pnma$ and $F$–$43m$ phases, there is no doubt that the Debye temperature of AlN is the largest, that of graphite is the second greatest, and that of cubic diamond is the lowest.

### Table 2. $v_s$, $v_p$, $v_m$ (m/s), and $\Theta_D$ (K) of XN in the $Pm$–$3n$, $Pnma$ and $F$–$43m$ phases.

| Material | Space Group | $\rho$ (g/cm$^3$) | $v_s$ (m/s) | $v_p$ (m/s) | $v_m$ (m/s) | $\Theta_D$ (K) |
|----------|-------------|-----------------|-------------|-------------|-------------|---------------|
| BN       | $Pm$–$3n$   | 2.829           | 9288        | 14749       | 10222       | 1571          |
| AlN      | $Pm$–$3n$   | 2.643           | 5604        | 9950        | 6235        | 793           |
| GaN      | $Pm$–$3n$   | 4.793           | 3794        | 6651        | 4217        | 515           |
| InN      | $Pm$–$3n$   | 5.291           | 2884        | 5389        | 3221        | 352           |
| BN       | $Pnma$     | 3.040           | 8642        | 14057       | 9537        | 1502          |
| AlN      | $Pnma$     | 2.828           | 5319        | 9508        | 5920        | 770           |
| GaN      | $Pnma$     | 5.114           | 3673        | 6633        | 4092        | 511           |
| InN      | $Pnma$     | 5.642           | 2595        | 5064        | 2907        | 325           |
| AlN      | $F$–$43m$   | 3.206           | 6169        | 10488       | 6837        | 927           |
| GaN      | $F$–$43m$   | 5.878           | 4226        | 7274        | 4690        | 613           |
| InN      | $F$–$43m$   | 6.496           | 2962        | 5493        | 3307        | 387           |

*Reference [16]. b Reference [18].

Mechanical anisotropy is often used to measure in which direction the maximum and minimum values exist. The three-dimensional (3D) graph can show this property intuitively, so the 3D maps of the $E$ for BN, AlN, GaN, and InN in the $Pm$–$3n$ phase are shown in Figure 4. If the mechanical properties of a material show isotropy in 3D space, the 3D silhouette maps of its mechanical properties are a regular sphere. If the shape is not a sphere, the mechanical properties of the material are anisotropic. At the same time, the less the shape of the 3D graph looks like a sphere, the greater the anisotropy [46]. Therefore, from the three-dimensional contour plots in Figure 4, it can be clearly seen that the mechanical anisotropy of the $E$ of AlN is the largest. Similar to other materials [47–50], the $Y_{max}/Y_{min}$ ratio (where $Y$ is $E$, $G$ and $v$) is used to quantify the anisotropy of various elastic moduli in this work. The maximum values and the minimum values of the $E$ for XN in the $Pm$–$3n$ phase are illustrated in Figure 5a, respectively. For Figure 5a, the blue and orange colours represent the $E_{max}$ and $E_{min}$, respectively. As shown in Figure 3b, the $E_{max}$ and $E_{min}$ of boron nitride in the $Pm$–$3n$ phase are still the largest, while that of indium nitride in the $Pm$–$3n$ phase is the lowest. The value of $E_{max}/E_{min}$ for $Pm$–$3n$ BN is shown in Figure 5d, and the blue, orange and grey colours represent the $E_{max}/E_{min}$ ratio, $G_{max}/G_{min}$ ratio, and $v_{max}/v_{min}$ ratios, respectively. It can be concluded from Figure 5d that the $E$ of $Pm$–$3n$ AlN has the largest anisotropy, as shown in Figure 4. Interestingly, in the $Pnma$ phase, the anisotropy of the $E$ of AlN is also the greatest [18], while the anisotropy of the $E$ of BN in the $Pm$–$3n$ phase is also the smallest, and the smallest anisotropy of the $E$ in the $Pnma$ phase is InN.
that of graphite is the second greatest, and that of cubic diamond is the lowest. This means that its Young's modulus is also the largest, the bond strength between the atoms is the largest, and the Debye temperature increases as the Young's modulus increases. The Debye temperature of XN (in the Pnma phase) is higher than that of BN (in the Pnma phase), typifies the number of atoms in the molecule, M. The crystal density, ρ, typifies the molecular weight, M. The bond strength between the atoms is the largest, which is reflected by the bulk modulus, B.

\[ B = \frac{1}{\rho} \left( \frac{\pi}{4} \rho \right)^{1/3} \]

Knowledge of B means that its value is a powerful tool to reflect the bonding force between atoms. The value of B typifies the bonding force between atoms. The value of B represents Boltzmann's constant, and D represents the bulk modulus and shear modulus of AlN, GaN and InN.

Table 2 shows the Young's modulus for BN, AlN, GaN, and InN. Figure 4 shows the three-dimensional contour plots of the Young's modulus for BN, AlN, GaN, and InN in the Pnma phase. Comparing the Young's modulus, as shown in Table 1, the three-dimensional contour plots for BN, AlN, GaN, and InN are shown in Figure 4.

Figure 4. The three-dimensional contour plots of the Young's modulus for BN (a), AlN (b), GaN (c), and InN (d) in the Pm-3n phase.

Figure 5. Cont.
To better and more easily comprehend the anisotropy of $E$, and the distribution of $E$ on the major planes of XN in the Pm–3n phase, such as the (001), (010), (100), (101), (110), (111), and (011) planes, the related $E_{\text{max}}$ and $E_{\text{min}}$ of the major planes are listed in Table 3. From Table 3 and Figure 5a, all the maximum values of $E$ for XN in the Pm–3n phase appeared in the (001), (010), (100), (110), and (011) planes, while the minimum values of $E$ of XN in the Pm–3n phase appeared in the (101), (110), and (011) planes. For the (111) plane, both the maximum values and minimum values of $E$ of XN in the Pm–3n phase do not appear in the (111) plane, the $E_{\text{max}}$ and $E_{\text{min}}$ for XN in the Pm–3n phase in the (111) plane are the same, and the $E_{\text{max}}$ and $E_{\text{min}}$ are the highest and the lowest values in all directions, which are shown in Figure 5a, so the values of $E$ of XN in the Pm–3n phase in the (111) planes are isotropic. Compared with BN in the Pnma phase, the anisotropy of $E$ along the (001), (010), (100), and (111) directions of BN in the Pm–3n phase is also significantly smaller than that in the Pnma phase. In addition, for the anisotropy of $E$ in the (001), (010), and (100) planes of AlN, GaN and InN in the Pm–3n and Pnma phases, the anisotropy of $E$ of the Pnma phase is almost the largest, except for the anisotropy of $E$ along the (100) direction of AlN in the Pm–3n phase, which is 1.79, and is slightly larger than that of the Pnma phase (1.77 [18]).

According to Hooke’s law and the ELAM software package [51], we know that $E$ can be expressed in two kinds of space angles, while the $G$ and $\nu$ need three kinds of space angle. The specific details
are described in references [51–53]. Therefore, the three-dimensional contour maps of the G and v are divided into 3D silhouette maps of the maximum values and minimum values. The 3D silhouette maps of the $G_{\text{max}}$, $G_{\text{min}}$, $v_{\text{max}}$ and $v_{\text{min}}$ for BN, AlN, GaN, and InN in the $Pm\overline{3}n$ phase are plotted in Figures 6a–d and 7a–d, respectively. Here, the dashed purple external and the solid purple external typify the $G_{\text{max}}$ and $G_{\text{min}}$, respectively, and the dashed orange external and the solid orange external typify the $v_{\text{max}}$ and $v_{\text{min}}$, respectively. The $G_{\text{max}}$, $G_{\text{min}}$, $v_{\text{max}}$ and $v_{\text{min}}$ and the $X_{\text{max}}/X_{\text{min}}$ ratio for XN in the $Pm\overline{3}n$ phase are plotted in Figure 5b–d. The blue colours typify the $G_{\text{max}}$ and $v_{\text{max}}$, and the orange colours typify the $G_{\text{min}}$ and $v_{\text{min}}$ in Figure 5b–c, respectively. The blue, orange and grey colours represent the $E_{\text{max}}/E_{\text{min}}$ ratio, $G_{\text{max}}/G_{\text{min}}$ ratio and $v_{\text{max}}/v_{\text{min}}$ ratio in Figure 5d, respectively. From Figures 6 and 7, it can be seen that the G and v of XN in the $Pm\overline{3}n$ phase exhibit elastic anisotropy, and from Figure 5b–d, AlN in the $Pm\overline{3}n$ phase has the greatest elastic anisotropy in G and v, BN in the $Pm\overline{3}n$ phase has the lowest elastic anisotropy in the shear modulus, and $Pm\overline{3}n$ InN has the smallest elastic anisotropy in the Poisson’s ratio. The distribution of G and v in the (001), (010), (100), (101), (110), (111), and (011) planes of XN in the $Pm\overline{3}n$ phase are also listed in Table 3. Different from the Young’s modulus, all the $G_{\text{max}}$, $G_{\text{min}}$, $v_{\text{max}}$ and $v_{\text{min}}$ of XN in the $Pm\overline{3}n$ phase appeared in the (001), (010), (100), (101), (110), (111), and (011) planes. In other words, the anisotropy of G and v of these seven main planes is the same. In addition, for BN in the $Pm\overline{3}n$ and $Pnma$ phases, both the elastic anisotropy in G and v in all directions of the $Pm\overline{3}n$ phase are smaller than that of the $Pnma$ phase. Furthermore, both the elastic anisotropy in G and v along the (001), (010), (100), and (011) directions of the $Pm\overline{3}n$ phase are weaker than that of the $Pnma$ phase [16,18]. Especially for v, the $v_{\text{max}}/v_{\text{min}}$ of the $Pm\overline{3}n$ phase is only 3.22, while the $v_{\text{max}}/v_{\text{min}}$ ratio of the $Pnma$ phase is 4.945, and the maximum value is in the (100) plane, which is as high as 14.431. The anisotropy of G and v of the $Pm\overline{3}n$ phase is obviously smaller than that of the $Pnma$ phase.

Figure 6. The 3D silhouette maps of the G for BN (a), AlN (b), GaN (c), and InN (d) in the $Pm\overline{3}n$ phase.
3.4. Electronic Properties

The electronic band structures of XN in the Pm-3n phase within the HSE06 function are illustrated in Figure 8a–d. The coordinates of the high symmetry points in the whole Brillouin zone for XN in the Pm-3n phase are X (0.500, 0.000, 0.000) → R (0.500, 0.500, 0.500) → M (0.500, 0.500, 0.000) → R (0.500, 0.500, 0.500). From Figure 8a–d, AlN, GaN, InN and BN can all be used as a semiconductor material, while AlN, GaN, InN are a direct band gap, and BN is an indirect and broad band gap. Compared with the Pnma phase, the band gap of BN in the Pm-3n phase is slightly less than that of the Pnma phase (7.18 eV within the HSE06 hybrid functional) [29]. For AlN, GaN and InN in the Pm-3n phase and Pnma phase with a direct band gap, the band gaps of the same compound in different phases are different. The band gaps of BN, AlN and GaN in the Pm-3n phase are slightly less than those of BN, AlN and GaN in the Pnma phase [16,18], while the band gap of InN in the Pm-3n phase is 1.04 eV, which is slightly larger than that of the Pnma phase [18]. In addition, the Fermi levels of XN in the Pm-3n and Pnma phases are also estimated in our paper. The related data of XN in the Pm-3n and Pnma phases are shown in Figure 8e. The Fermi levels of XN in the Pm-3n and Pnma phases decrease in turn; the difference is that the reduction degrees of the Pm-3n phase and Pnma phase are different. The Fermi level of XN in the Pm-3n phase decreases by 7.86 eV, and the reduction degree of XN in the Pnma phase is smaller than that of the Pm-3n phase.

Figure 7. The 3D silhouette maps of the v for BN (a), AlN (b), GaN (c), and InN (d) in the Pm-3n phase.
4. Conclusions

According to DFT, the structural properties, stability, elastic properties, elastic anisotropy properties and electronic performances of BN, AlN, GaN and InN in the $Pm\overline{3}n$ phase are estimated in our paper. The lattice parameters of InN in the $Pm\overline{3}n$ phase increase by 40.54% compared to those of BN in the $Pm\overline{3}n$ phase, and this increase is larger than that of the lattice parameter $a$ of InN in the $Pnma$ phase to BN in the $Pnma$ phase ($a$: 38.59%). This increase is slightly lower than that of lattice parameters $b$ and $c$ of InN in the $Pnma$ phase to BN in the $Pnma$ phase ($b$: 42.87% and $c$: 41.26%). $Pm\overline{3}n$ BN is the most brittle, and $Pm\overline{3}n$ InN is the most ductile. $Pm\overline{3}n$ GaN is between the brittleness and ductility of $Pm\overline{3}n$ BN and $Pm\overline{3}n$ InN but tends to be brittle. The calculated Debye temperature of XN in $Pm\overline{3}n$ phase follows the order InN $<$ GaN $<$ AlN $<$ BN. By showing the three-dimensional
contour plots and the $Y_{\text{max}}/Y_{\text{min}}$ ratio (where $Y$ is shear modulus $G$, Young’s modulus $E$, and Poisson’s ratio $\nu$), the anisotropy of $E$, $G$, and $\nu$ of the $Pm\bar{3}n$ phase is obviously lower than that of the Pnma phase. The electronic band structures of XN in $Pm\bar{3}n$ phase show that AlN, GaN and InN are semiconductor materials with direct band gaps within the HSE06 hybrid functional, while BN in the $Pm\bar{3}n$ phase is an indirect and wide semiconductor material. Compared with III-V nitrides compounds in $Pnma$ phase, the band adjustable range in $Pm\bar{3}n$ phase (1.04–5.87 eV) is slightly smaller than $Pnma$ phase (0.66–7.18 eV). It is likely to have a good adhibition in the electronic manufacturing industry such as LEDs, UV detectors, infrared detectors and visible light detectors. Compared with the previous materials (GaAsP, AlGaN, AlGaAs, AlGaN and other ternary semiconductor alloys, or AlGaNp quaternary semiconductor alloy), they save the trouble of making ternary or quaternary semiconductors in semiconductor technology.

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