APOGEE KINEMATICS. I. OVERVIEW OF THE KINEMATICS OF THE GALACTIC BULGE AS MAPPED BY APOGEE

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Received 2015 September 25; accepted 2015 December 14; published 2016 February 23

ABSTRACT

We present the stellar kinematics across the Galactic bulge and into the disk at positive longitudes from the SDSS-III APOGEE spectroscopic survey of the Milky Way. APOGEE includes extensive coverage of the stellar populations of the bulge along the midplane and near-plane regions. From these data, we have produced kinematic maps of 10,000 stars across longitudes of 0° < l < 65°, and primarily across latitudes of |b| < 5° in the bulge region. The APOGEE data reveal that the bulge is cylindrically rotating across all latitudes and is kinematically hottest at the very center of the bulge, with the smallest gradients in both kinematic and chemical space inside the innermost region (l, |b|) < (5°, 5°). The results from APOGEE show good agreement with data from other surveys at higher latitudes and a remarkable similarity to the rotation and dispersion maps of barred galaxies viewed edge-on. The thin bar that is reported to be present in the inner disk within a narrow latitude range of |b| < 2° appears to have a corresponding signature in [Fe/H] and [α/Fe]. Stars with [Fe/H] > −0.5 have dispersion and rotation profiles that are similar to that of N-body models of boxy/peanut bulges. There is a smooth kinematic transition from the thin bar and boxy bulge (l, |b|) < (15°, 12°) out to the disk for stars with [Fe/H] > −1.0, and the chemodynamics across (l, b) suggests that the stars in the inner Galaxy with [Fe/H] > −1.0 originate in the disk.

Key words: stars: abundances – stars: fundamental parameters – stars: kinematics and dynamics – surveys – techniques: spectroscopic

Supporting material: machine-readable table

1. INTRODUCTION

The Galactic bulge of the Milky Way is a cylindrically rotating system (e.g., Howard et al. 2008; Kunder et al. 2012; Zoccali et al. 2014) for the stars with [Fe/H] > −1.0 (Ness et al. 2013b). These [Fe/H] > −1.0 stars, which compose the majority of stars in the bulge (e.g., Zoccali et al. 2008), have a smooth kinematic transition out to the disk, as well as symmetry about the major axis. Previous studies of the bulge, which have mostly examined stars at |b| > 4°, have demonstrated that there is a kinematic dependence of the stars in the bulge on [Fe/H]. The stars with [Fe/H] > −0.5, which are part of the X-shaped profile of the Milky Way (Ness et al. 2012; Utenthaler et al. 2012) that was found in the red clump density distribution (Nataf et al. 2010; McWilliam & Zoccali 2010), show a velocity dispersion profile that is consistent with bulge formation via dynamical instabilities seen in N-body models. The most metal-rich subset of these stars (i.e., [Fe/H] > 0) are kinematically colder than their more metal-poor counterparts (−0.5 < [Fe/H] < 0), with the same characteristic dispersion profile, as a function of (l, b). This dependence of the kinematic profile on metallicity for stars that appear to be part of the bulge/bar is not yet fully understood, but is likely related to the mapping of the stars into the boxy bulge from the disk by the bar, as a function of their initial phase space in the disk (Martínez-Valpuesta & Gerhard 2013; Di Matteo et al. 2015).

Bulge stars in the lower-metallicity range of −1.0 < [Fe/H] < −0.5 show a similar rotation, but dissimilar dispersion, to the more metal-rich stars. The dispersion profile of these more metal-poor stars does not display a strong dependence on latitude and appears dissimilar to that of N-body models of bulge formation via dynamical instabilities of the bar. These stars are possibly part of the thick-disk population in the inner Milky Way, given their fast rotation and high, latitude-independent dispersion (e.g., Ness et al. 2013b; Di Matteo et al. 2015). Alternatively, these stars may be part of a unique, classical bulge component of the Milky Way (e.g., Babusiaux et al. 2010; Hill et al. 2011). It has recently also been suggested that these stars could also have an origin in massive clumps at higher redshift (Zoccali et al. 2014). The RR Lyrae stars, which peak at a metallicity of [Fe/H] < −1.0 in the bulge, follow an ellipsoidal density distribution along the
The bulge fields of the APOGEE survey (Majewski et al. 2012; Majewski et al. 2015) are observed in the infrared and thus are uniquely and critically placed to characterize previously unexamined regions of the bulge. Many previous bulge studies have revealed the kinematic trends for fields in the bulge located off the plane (|b| > 4°). Additionally, the GIBS survey (Zoccali et al. 2014) has examined 14 fields within (l, b) < (10°, 4°), as part of their combined 24-field sample of 5000 red clump giants (plus an additional 12000 red clump giants from archive data) toward the bulge. Using the APOGEE data, which trace unexamined regions across the bulge into the disk, including at b < 4°, we address the following questions: (1) What is the nature of the kinematic profile of the bulge in previously unexplored regions in and near the plane? (2) Are the kinematics in these previously unexplored regions at odds or consistent with the kinematics measured off the plane and with N-body models? (3) What is the dependence of the kinematic profile of the stars on [Fe/H]?

The high-resolution, 22,500 spectra spanning 1500–1700 nm allow not only the first homogeneous, contiguous mapping of the bulge metallicity and kinematics through the plane of the Milky Way into the innermost bulge region but also a mapping of the kinematics as a function of multielement abundances. For the characterization of stellar populations in these previously unexplored regions presented in this paper, we use the APOGEE spectra from the SDSS-III Data Release 12 (DR12; Ahn et al. 2014) for about 20,000 stars toward the Galactic bulge and surrounding disk (Section 2). These stars span mostly positive longitudes, from −5° < l < 65°, across |b| < 5° and including some higher-latitude fields at b < 15°. We have implemented distance cuts in order to select the 10,000 stars that are located at distances 4–12 kpc along the line of sight.

2. STELLAR PARAMETERS AND DISTANCES

The APOGEE survey, part of the SDSS-III project (Eisenstein et al. 2011), operates at the 2.5m telescope of the Apache Point Observatory (Gunn et al. 2006). Being a near-IR survey, APOGEE has observed a large, homogeneous sample of stars (~150,000) spanning the halo, disk, and innermost regions of the Milky Way, including the boxy bulge (for details of the target selection, see Zasowski et al. 2013). All data are processed with a custom pipeline (Nidever et al. 2015), which also calculates the stellar radial velocities with an accuracy of <1 km s\(^{-1}\).

We use the publicly available spectra, released as part of DR12, to determine stellar parameters of \(T_{\text{eff}}, \log g, [\text{Fe/H}], \) and \([\alpha/\text{Fe}]\) using The Cannon (Ness et al. 2015). The set of reference objects for training The Cannon is composed of a subsample of ~2150 high signal-to-noise ratio (S/N ~ 300) APOGEE stars with no STAR BAD set in ASPCAPFLAG, as well as the globular and open cluster stars (Mészáros et al. 2013). We adopt the high-S/N APOGEE data, as described in Ness et al. (2015). We use the corrected \(T_{\text{eff}}, \log g, [\text{Fe/H}], \) and \([\alpha/\text{Fe}]\) labels from ASPCAP (García Pérez et al. 2015) for these training stars (the corrections to these parameters are done by APOGEE using the cluster calibration sample; Mészáros et al. 2013). We also include the metal-poor ([Fe/H] < −1.0) globular cluster stars with their isochrone-corrected training labels (Ness et al. 2015) so as to extend our training set to the lowest metallicities. The Pleiades open cluster, which contains the main-sequence stars in the training set, enables dwarf stars in the DR12 data to be identified by The Cannon and therefore eliminated. In total, there are 2400 reference objects that are used for the training step of The Cannon. The \([\alpha/\text{Fe}]\) from the APOGEE spectra is composed of contributions from many alpha-elements, including O, Ca, Ti, Si, and Mg (see García Pérez et al. 2015).

For our analysis, standard and telluric stars were removed from the sample, as well as flagged rapid rotators and special targets from ancillary programs. The main survey objects were selected using the APOGEE EXTRATARG flag. Only giants were chosen for the analysis with \(\log g < 3.8\). Following this selection, there remain ~20,000 stars in the direction of the bulge and inner disk, across |l| < 65° and |b| < 15°. These ~20,000 stars are shown in Figure 1 (no distance cuts have yet been applied to this sample). We note that these have had no corrections or calibrations applied to them, nor prior information about the \(T_{\text{eff}}–\log g\) space of the isochrones; these results are directly out of The Cannon’s label-estimation step.

The parameter range of The Cannon’s training set is narrower than that of the APOGEE stars toward the bulge and disk. The training set spans ranges of 3600 K < \(T_{\text{eff}} < 5750\) K, −3 < \(\log g < 4.7\), −2.5 < [Fe/H] < 0.45, and −0.3 < [\(\alpha/\text{Fe}\)] < 0.5. The range of parameters of the giant stars toward the bulge and disk determined by The Cannon for the survey spectra is 3500 K
Table 1
Partial Column Excerpt from the Online Table of Four Stellar Labels (Teff, log g, [Fe/H], [α/Fe]) Determined by The Cannon for APOGEE DR12 Stars toward the Bulge and Inner Disk with Their $\chi^2_{\text{reduced}}$ Values and Distances

| Star ID (2MASS) | Teff (K) | log g (dex) | [Fe/H] (dex) | [α/Fe] (dex) | $\chi^2_{\text{reduced}}$ | Distance (kpc) | $V_G$ (km s$^{-1}$) |
|-----------------|----------|-------------|--------------|--------------|----------------|---------------|-------------------|
| 2M18350759-0508445 | 3727.8 | 0.7 | -0.18 | -0.0 | 1.5 | 12.1 | 159.1 |
| 2M18354096-0452009 | 3952.6 | 1.4 | 0.2 | -0.03 | 1.4 | 5.5 | 130.1 |
| 2M18354152-0533504 | 3952.8 | 0.7 | -0.34 | 0.04 | 3.1 | 9.7 | 157.4 |
| 2M18354345-0453021 | 4636.0 | 2.6 | 0.37 | -0.0 | 1.2 | 2.9 | 120.8 |
| 2M18354380-0425029 | 4053.5 | 1.4 | -0.44 | 0.23 | 1.2 | 10.6 | 205.7 |
| 2M18355233-0501398 | 4604.9 | 2.3 | 0.35 | 0.01 | 0.9 | 1.1 | 93.8 |
| 2M18361149-0608046 | 4418.3 | 1.9 | 0.34 | 0.02 | 1.3 | 1.6 | 128.6 |
| 2M18361409-0456347 | 4159.6 | 1.8 | 0.52 | -0.07 | 2.6 | 3.9 | 141.7 |
| 2M18362011-0510251 | 4432.5 | 2.1 | 0.21 | 0.03 | 3.5 | 3.8 | 145.8 |
| 2M18362181-0545341 | 4625.9 | 2.5 | -0.36 | 0.21 | 2.0 | 2.5 | 131.7 |

Note. APOGEE targeting flags should be used in the selection of stars, and main-sequence stars, are provided to identify these stars as such, but stellar parameters are unreliable for these stars, which should be excluded from analysis.

(This table is available in its entirety in machine-readable form.)

Figure 2. Median distance showing the 1$\sigma$ dispersion in the error bars, of the APOGEE stars as a function of longitude, for |b| < 2° (left), 2° < |b| < 4° (center) and 4° < |b| < 8° (right).

$< T_{\text{eff}} < 5500$ K, -0.5 < log $g$ < 3.8, -2.0 < [Fe/H] < 0.7, and -0.3 < [α/Fe] < 0.5. APOGEE stars that have stellar parameters determined by The Cannon outside of that of the training data are assigned values based on extrapolation of The Cannon’s model. Therefore, the scale of these parameters is not tied directly to any physical calibration.

Nevertheless, as seen in Figure 1, almost all stars fall naturally on the isochrone of the appropriate metallicity. The parameter space of the survey stars that extend beyond that of the training data extrapolates directly along the isochrones to the lowest log $g$ and Teff values. The stellar parameters returned by The Cannon also compare well with the results from ASPCAP, where the ranges overlap. For the ≈18,000 stars in the range (-5° < l < 65°, |b| < 15°) for which APOGEE stellar parameters are available from ASPCAP, The Cannon’s parameters are identical within $\Delta$[Fe/H] = -0.04 ± 0.07, $\Delta$[α/Fe] = 0.01 ± 0.06, $\Delta$Teff = -22 ± 55 K, and $\Delta$log $g$ = 0.02 ± 0.12.

The stellar parameters for these stars, their $\chi^2_{\text{reduced}}$ values, and their corresponding distances, are provided in Table 1. Stars with a poor fit to the model $\chi^2_{\text{reduced}} > 6$ were excluded from analysis (approximately 5% of the stars).

We do include the stars in this region of extrapolated stellar parameter space in our analysis of the kinematics of stars as a function of metallicity. However, we use only coarse bins in metallicity of 0.5 dex intervals (see Section 5). This allows us to compare the kinematic behavior of metal-rich and metal-poor stars. The alternatives, of either eliminating all stars in the region of extrapolated label space or using only the stars with stellar parameters provided by APOGEE (which are also in extrapolated label space for ASPCAP), produce qualitatively the same results as those obtained with their inclusion. However, we have the benefit of working with a larger number of stars using the full range of The Cannon’s results.

To isolate the stars in the region of the bulge, we use the stellar parameters and obtain distances by shifting the stars in Teff and log $g$ space to their nearest position on a 10 Gyr isochrone (Bressan et al. 2012), for an isochrone at the metallicity of the star by interpolating between isochrones. The apparent dereddened magnitude of each star was calculated by adopting the RJCE-WISE extinction value for that line of sight, provided in the APOGEE DR12 data (Majewski et al. 2011; Zasowski et al. 2013). In the few cases, for 160 stars, where no reddening estimates were available, a value of $A(K_s) = 0$ was adopted, which effectively provides an upper limit to the distance estimate from the isochrone matching. The adopted distance error on each star is 30%. These errors are composed of uncertainties in the reddening, which is most significant in the plane, uncertainties on the stellar parameters (which are $\Delta$ (Teff, log $g$, [Fe/H], [α/Fe]) = 70 K, 0.2, 0.1, 0.05 dex), and also any systematic offsets in the stellar parameters or else the isochrones.

Figure 2 shows the distance distribution of stars across longitude, colored by metallicity, for three slices in latitude. This shows that using distances calculated directly from the isochrones, the stars in the APOGEE sample are biased to be located on the near side of the bulge and there is a strong bias
in metallicity, with few metal-rich stars beyond $R_{GC} > 8$ kpc. In addition to the large uncertainties in distance for the red giant selection ($\approx 30\%$), particularly in highly extinguited regions, there may be systematic offsets in the distance and the median distance may be underestimated by this simple isochrone distance calculation. We examine the chemo-kinematics of stars well in front of the bulge ($<3$ kpc) and for our bulge selection (4–12 kpc) in Sections 3–5 to validate that the relative scale of our distances effectively enables the foreground stars, which show disk-like kinematics and chemistry that is dissimilar to the bulge sample, to be effectively eliminated from our analysis. Systematic distance offsets may be a consequence of the stellar parameter scale (e.g., temperature scale) of the stars or else overestimated values of extinction, particularly in the highly reddened plane region ($|b| < 4^\circ$). The metallicity bias as a function of distance in the APOGEE sample is discussed in A. E. Garcia-Perez et al. (2016, in preparation).

The bias in distance as a function of [Fe/H] and also the concentration in distance to the near side of the bulge are present even where the extinction is smaller (above $|b| > 4^\circ$). The peak of the overall distance distribution is similar to that found by Ness et al. (2013a) for red clump stars, at around $5–6$ kpc from the Sun although the APOGEE stars, and particularly the metal-rich stars do not reach far beyond $8$ kpc.

For our distance cuts in our selection of bulge stars, we adopt a broad heliocentric distance criterion of $4–12$ kpc for bulge membership, which includes stars from the end of the bar at $l = 30^\circ$ to the far side of the bar at $l < 0^\circ$. Although broad, this distance selection effectively and critically eliminates many of the stars in the foreground disk along the line of sight to the bulge while including the stars in the bulge, the vast majority of which are on the near side. Our distance estimates are included with the stellar parameters for DR12 in Table 1.

Our distance cut of $4–12$ kpc provides a sufficient margin to account for uncertainties in the distance determination, while minimizing the contamination of the foreground objects. Many previous bulge studies (e.g., Howard et al. 2008; Zoccali et al. 2008) do not determine distances to eliminate foreground objects; however, these studies do select stars in a confined region of the color–magnitude diagram to maximize the likely true bulge stars in the sample. We find with our determined distances and APOGEE color cuts (Zasowski et al. 2013) that approximately 35% of the sample consists of foreground objects along the line of sight toward the bulge located at distances $<4$ kpc (which we correspondingly eliminate; a similar fraction as found by Ness et al. 2013a). The distances in the plane show bias to distances nearer to the Sun, and the median distances increase above the plane ($|b| > 2^\circ$) in regions of lower extinction, where the fainter stars are located at larger distances.

Figure 3(a) shows the mean distance distribution across ($l, b$) with our distance cuts implemented; off the plane ($|b| > 2^\circ$), the mean distances are around $6.5–7$ kpc from the Sun, and Figure 3(b) shows that the distance distribution is broader at higher latitudes. Adopting instead the distance estimates of Hayden et al. (2015) for our stars and repeating the analyses below, we obtain no significant differences in our results and we draw no different conclusions from those in this analysis with The Cannon-derived parameters and our own distances.

After our distance cuts, our sample of bulge and inner disk stars is reduced to $\sim 10,000$. For these stars, we examine the mean velocity and dispersion by mapping across bins with $\Delta l = 2^\circ$ and $\Delta b = 1^\circ$ across the bulge and inner disk (Section 3). The typical number of stars in each of these bins is $\sim 20$ as shown in Figure 4.

3. KINEMATIC MAPS OF THE MILKY WAY

The rotation and dispersion maps of the bulge and inner disk, comprising the 10,000 stars spanning estimated distances of $4–12$ kpc, are presented in Figure 5. The color map shows the mean velocity and the dispersion in each bin. Figure 5(a) demonstrates that the rotation curve appears cylindrical into the plane, with about the same average velocity across slices in latitude. Figure 5(b) reveals that the bulge is kinematically hottest at the very center and the dispersion decreases with ($l, b$) and flattens as a function of $b$ at the disk $l > 15^\circ$. The comparable results from BRAVA (Howard et al. 2008; Kunder et al. 2012) and ARGOS (Freeman 2012; Ness et al. 2013b) are also included on these maps. The APOGEE velocity data agree well with the results from the BRAVA and ARGOS surveys, which show comparable mean rotation and dispersion values for fields that overlap and a near-cylindrical rotation profile across high to low latitudes. The rotation and dispersion profiles transition smoothly into the disk and do not reveal any large-scale dynamical substructure present in the bulge.

Overall, these maps are remarkably similar to the cylindrical rotation and velocity dispersion profile of barred galaxies (e.g., NGC 7332 by Falcón-Barroso et al. 2004) and also the corresponding maps of N-body simulations, which evolve boxy bulges from the disk (e.g., Athanassoula & Misiriotis 2002). Two critical differences between the IFU maps of external galaxies and our maps are that we have eliminated the foreground objects and that these averages are from a star-by-star analysis, rather than an integrated stellar population.

We examine the comparable kinematic profile for the 6800 foreground objects, which are shown in number distribution in Figure 6, where stars have been selected with distances from the Sun $<3$ kpc. The rotation of the disk stars nearer the Sun presented in Figure 7(a) is similar to the stars at the distance of the bulge, although the rotation is slower. However, the velocity dispersion shown in Figure 7(b) is clearly dissimilar to that of the bulge population, with a flatter, low-dispersion profile from inner to outer longitudes, as expected for a foreground disk population.

4. COMPARISON TO AN N-BODY MODEL

4.1. Rotation and Velocity Dispersion Profiles of the Bulge

The overall kinematic trends from the ARGOS survey, which comprised 28 fields of about 600 stars at distances of the bulge (4.5–11.5 kpc from the Sun), across $b = -5^\circ, -7.5^\circ$, and $-10^\circ$ and longitudes $|l| < 31^\circ$, are a good match to $N$-body models of boxy/peanut bulges for stars with metallicities of $[\text{Fe/H}] > -0.5$ (Ness et al. 2013b). In this section, we compare the global APOGEE kinematic patterns to a model of Athanassoula (2005, 2007), which has been scaled to match the ARGOS data in Ness et al. (2013b).

This model consists of a disk and dark matter halo with density distributions as described in Athanassoula (2007). The model is evolved from a thin disk that forms a bar at early times, whose instabilities later form a boxy/peanut bulge. Because the model contains no gas, there is no star formation and there is no information on chemistry of the stars. The model has a Toomre parameter of $Q = 1.2$ and a vertical height
Figure 3. (a) Mean distance and (b) distance dispersion maps for the 10,000 APOGEE bulge and disk stars, isolated with a heliocentric distance cut of 4–12 kpc. Note that the stars in the plane are nearer to the Sun and that the range of distances spanned at higher latitudes is largest above the plane.

Figure 4. Number of stars in each bin used for the determination of the mean velocity and velocity dispersion maps, for stars with distances of 4–12 kpc. There are a total of ≈10,000 stars at these distances for these fields.
Figure 5. (a) Rotation and (b) velocity dispersion maps for the 10,000 APOGEE bulge and disk stars, taken within a distance cut of 4–12 kpc. The squares with black outlines are from the ARGOS data and contain about 600 stars. The circles with black outlines are BRAVA fields and contain between 100 and 200 stars in each of these fields. The number of stars in each of the APOGEE fields is given in Figure 4. Typical errors in the velocities for individual APOGEE stars are <1 km s\(^{-1}\). These maps show a cylindrical rotation that is symmetrical about the major axis and a dispersion profile that is hottest at the center, similarly to the kinematic profiles of boxy bulges of other galaxies. There is also good agreement between APOGEE, ARGOS, and BRAVA results.

Figure 6. Numbers of stars in each bin used for the determination of the mean velocity and velocity dispersion maps, for stars with distances of 0–3 kpc. There are a total of ≈6800 stars at these distances for these fields.
that is 0.2 times the disk scale length. This model is scaled to the spatial size of the COBE bar. Ness et al. (2013b) verified this model to be a good match to the ARGOS kinematics at latitudes below the plane, from $b = -5^\circ$ to $-10^\circ$.

This model was compared in Ness et al. (2013b) to the ARGOS data across latitudes of $b = -5^\circ$, $-7.5^\circ$, and $-10^\circ$ and found to be a good qualitative match to the data. The APOGEE data have stars across a broader spatial region of the bulge and, critically, in the plane. Figure 8 shows the comparison of the model with the ARGOS fields across the three latitudes plus the APOGEE data in the plane, where APOGEE stars that are within $b < 2^\circ$ have been binned to compose the in-plane measurements. The bar in the model has been rotated to $27^\circ$, and stars in the model have been selected at distances 4–12 kpc from the Sun, similar to the APOGEE and ARGOS data.

There are small quantitative differences between data and the model seen in Figure 8, but the trends are matched very well by the model. Particularly, inside the boxy bulge region, the dispersion is peaked at low latitudes and latitude dependent. At $b = 0^\circ$, the dispersion in the plane is almost double the dispersion at $b = 10^\circ$. Outside of $l > 10^\circ$, however, the disk shows comparable dispersion across latitude, at $b = 0^\circ$ and $b = 10^\circ$.

4.2. The Velocity and Velocity Dispersion Maps of the Bulge

Mean velocity and velocity dispersion maps, comparing data and the model, for APOGEE, BRAVA, and ARGOS, across the bulge and moving into the disk, are shown in Figures 9 and 10, respectively. These figures include maps of the residuals between data and the model and the errors on these residuals, for the ARGOS and APOGEE fields. Panels (a)–(d) include stars only within distances 4–12 kpc from the Sun (except for the BRAVA data, where no stellar parameters or subsequently distances were determined for the stars). Panels (a) and (b) show the rotation of the data and model and the dispersion of the data and model, in Figures 9 and 10, respectively. Panels (c) and (d) then show the difference between the model and data (model–data) and the error on that value, respectively.

From Figure 9, panels (a) and (b), it is apparent that the overall mean velocity of the model is a good match to the data. The residuals in panel (c) do not show coherent angular structure. There are a few mismatching bins between data and model but with correspondingly large errors due to the small uncertainties.
numbers of stars along these lines of sight (see panel (d)). Note in Figure 9(d) that the errors are relatively low for the ARGOS fields as there are about 600 stars that compose these fields, in comparison to 30 on average (but in some cases lower) in the binned APOGEE fields composing the rest of the map.

In the dispersion profiles shown in Figure 10, and as also seen in the bottom panel of Figure 8, this particular model is slightly kinematically hotter than these data, within |l| < 5°. As noted by Ness et al. (2013b), the model is more centrally concentrated than the data. With the exception of the few bins in the midplane distributed across a narrow range in latitude, there is no overall structure in the difference in dispersion between model and data, as seen for the rotation. Some bins in the midplane are, however, significantly kinematically cooler than the model. This is seen just below the plane, for a narrow latitude range (b ∼ 1°), and this extends from l = 12° out into the disk. There are very few stars along the line of sight that are located within this apparently kinematically cool structure in the disk, although the sampling errors shown in Figure 10(d) are small and indicate that these differences are statistically significant.

Figure 11 shows the numbers of stars in the fields in Figures 9(a) and 10(a). Figure 2 shows that the stars within |b| < 2° are relatively nearer to the Sun than compared to stars farther from the plane. Although Figure 10(d) indicates that the low dispersion in some bins near the plane in the data compared to the model is statistically significant, this lower velocity dispersion that is measured may simply reflect the limited extent in distance of these observed stars compared to the surrounding stars. The narrow distance dispersion of the stars may also be a consequence of the stellar density, which is expected to peak at the bar (which is located at about 4.7 kpc away at 30°). It is possible that this kinematically cool component is associated with the “thin” bar or “super-thin” bar structure reported by Wegg et al. (2015; see Section 5.1). These data may alternatively reflect the presence of a very cool part of the disk that extends from the outer region into the bulge (rather than being a “thin” bar).

Overall, the kinematic behavior of the APOGEE bulge stars is in good agreement with the behavior in the previous fields mapped by ARGOS and BRAVA and with N-body models of boxy/peanut bulges formed from bars via instabilities of the Galactic disk (see also Shen et al. 2010; Portail et al. 2015).

5. CHEMODYNAMICAL SIGNATURES OF THE BULGE COMPONENTS

The kinematics of stars in the bulge is a function of the [Fe/H], and the rotation and dispersion profiles of the stars change with metallicity (Babusiaux et al. 2010; Ness et al. 2013b). The APOGEE data allow for the examination of the kinematics not only as a function of [Fe/H] for the previously unexamined midplane and low-latitude regions but also for [α/Fe] and [X/Fe]. In this first chemodynamical analysis, we report the global trends with [Fe/H].
5.1. The Long Bar

The boxy/peanut extent of the bulge, detected in the COBE/DIRBE image (Dwek et al. 1995), extends to a latitude above the plane of around \( b \approx 12^\circ \) and a longitude (on the near side) of about \( l \approx 15^\circ \). The boxy/peanut region of the bulge grows out of a thinner bar—in the innermost boxy region, the stars that were originally in this bar are on orbits now extended in the vertical direction (e.g., Athanassoula 2005; Debattista et al. 2005; Martínez-Valpuesta et al. 2006).

As expected (Athanassoula 2005), this thinner part of the bar reaches further in the plane of the Galaxy than the \( l = 15^\circ \) boxy extent. Wegg et al. (2015) examined the red clump density distribution of stars toward the bulge using VVV photometry and report a \( |z| \lesssim 180 \text{ pc} \) thin bar and \( |z| \lesssim 45 \text{ pc} \) super-thin bar. The latter of these extends from the center of the bulge to longitudes \( l \approx 25^\circ \)–\( 30^\circ \). The thinnest component has the highest density at the end of the bar, \( l \sim 30^\circ \). The transition between the super-thin bar and thin bar structures occurs at a latitude of \( |b| \approx 2^\circ \). Figure 10 demonstrated the presence of a kinematic structure not seen in the \( N \)-body model, which, given its limited extent in \( b \), is more likely associated with the super-thin bar than the thin bar, but may also be simply part of the disk that is being observed across a narrow range in distance. We now look for signatures of this thin bar in the APOGEE data that might correlate with the morphology, \([\text{Fe/H}], \) and \([\alpha/\text{Fe}]\).

The mean \([\text{Fe/H}]\) and \([\alpha/\text{Fe}]\) of stars in the bulge and inner disk are shown in Figures 12 and 13. The trends revealed in these maps are discussed in detail in A. E. García-Pérez et al. (2016, in preparation). These trends are sensitive to the derived distances and subsequent distance cuts and selection of stars included in the analysis. Care must be taken comparing the mean metallicity across longitude, as the APOGEE target selection is different for bulge (357° \( \leq l \leq 22^\circ \)) and disk (see Zasowski et al. 2013). Brighter stars are targeted in the bulge, and the bulge metallicity distribution contains more cooler, lower \( \log g \) stars, compared to the disk sample. Figures 12(a) and 13(a) show the \([\text{Fe/H}]\) and \([\alpha/\text{Fe}]\) maps made for all 10,000 stars at distances of 4–12 kpc. The majority of foreground stars, which are shown in Figure 15 for a distance selection of \(<3 \text{kpc}\), are metal-rich stars; these metal-rich foreground stars are eliminated from the bulge maps by performing these distance cuts.

For comparison, Figures 12(b) and 13(b) show the sample of these stars with the same distance cuts, but with a lower \( \log g \) limit imposed of \( \log g > 0.5 \text{ dex} \). This removes a large number of the lowest \( \log g \) stars, which are targeted in a larger fraction in the bulge compared to the disk, owing to the brighter magnitude limit of the bulge selection strategy. This effectively eliminates many of the more distant, low-metallicity stars that have higher relative fractions within \( l < 22^\circ \). This demonstrates how the metallicity map across longitude is sensitive to distance biases. Consistent qualitative trends, however,
illustrated in Figures 12(a) and (b) and 13(a) and (b), are that the highest-metallicity stars, with corresponding lowest [α/Fe] enhancements, are concentrated to within |b| < 2°. The apparent metal-rich structure in the midplane (b < 2°), with a corresponding low α-enhancement (as seen in Figure 13), may be associated with the thin bar of Wegg et al. (2015).

Figures 12 and 13 show that there is a homogeneity in the mean [Fe/H] and [α/Fe] within (l, b) < (5°, 5°), the innermost region of the boxy part of the bar. There is a known metallicity gradient with latitude in the bulge, for latitudes |b| > 5°, of about −0.45 dex kpc⁻¹ (e.g., Zoccali et al. 2008; Ness et al. 2013a; Rojas-Arriagada et al. 2014) and a smaller gradient in longitude (e.g., Gonzalez et al. 2013). The flattening of the metallicity gradient with latitude near the plane has been detected (Ramírez et al. 2000; Rich et al. 2007, 2012), and the APOGEE survey now demonstrates that flattening is present not only at the minor axis but also at the entire extent of the boxy bulge (l < 15°) within |b| < 2°.

There is a relatively sudden decrease in [Fe/H] at |b| > 5° above the plane, seen in Figures 12 and 13 across l. There is also a sudden increase in [α/Fe] for all latitudes above b > 2° as seen in Figure 13. The metal-rich stars concentrated in the plane bar may reflect that there is ongoing star formation and even a potentially younger population in the plane extending out into the disk (Ness et al. 2014).

5.2. The Long Bar and Boxy Bulge

Figure 14 shows the dispersion in [Fe/H] and [α/Fe] of the bulge and disk. These trends in [Fe/H] and [α/Fe] dispersion are very similar for both selections compared in Figures 12(a), (b) and 13(a), (b), so only the full selection in (a) of these figures is shown. Figure 14 illustrates that the MDF is broadest at the minor axis. As seen in Figure 14(a), the [Fe/H] dispersion is largest on the minor axis, at all b, and as seen in Figure 14(b), the [α/Fe] dispersion is highest within (l, b) < (5°, 5°). The broadest MDF in the innermost region may reflect the mixture of populations in the center of the Milky Way (e.g., Ness et al. 2013a; Gonzalez et al. 2015). The population at high latitudes is metal-poor and alpha-rich (see Figures 12 and 13), and, at high latitudes, both disk and bulge have very similar mean metallicity, ⟨[Fe/H], [α/Fe],⟩ and metallicity dispersion, σ([Fe/H], [α/Fe]). This suggests a homogeneous, likely thick-disk population extending to the innermost region of the Milky Way, similar to the thicker part of the disk seen near the Sun (Ness et al. 2013a). Correspondingly, the kinematic behavior is smooth moving from the inner to the outer region, as seen from Figure 5.

The metallicity dispersion in the bulge shown in Figure 14 is >50% lower outside of l > 10° compared to inside of l < 10°. The lowest dispersion in [Fe/H] is in the plane in the disk at l > 30°, within |b| < 2° (Figure 14(a)). The dispersion in [α/Fe] decreases by about 50% outside of (l, b) < (10°, 5°) (Figure 14(b)) and also decreases above b > 10° at the minor axis.

Figure 5(a) shows that the overall trends in the velocity are smooth, with no obvious line-of-sight velocity signatures associated with the changing gradients in metallicity across (l, b) seen in Figure 12. In Figure 5(b), we see that the bulge is kinematically hottest in the innermost region (l, b) < (5°, 5°), and the dispersion drops off rapidly in b outside of (l, b) < (5°, 5°); this pattern corresponds to where the mean metallicity drops in
latitude, for longitudes within $l < 5\,\degree$. Otherwise, at lower latitudes, the metallicity within $l < 5\,\degree$ is fairly constant. Outside of longitudes $l > 10\,\degree$, the dispersion, shown in Figure 5(b), decreases gradually into the disk and becomes nearly constant with latitude as longitude increases.

Figure 15 contrasts the [Fe/H] and [$\alpha$/Fe] map of stars in the bulge with the foreground disk stars with distances $<3\,\text{kpc}$ from the Sun. The foreground disk is homogeneously metal-rich to latitudes $b < 10\,\degree$. It is also alpha-poor at low latitudes, $b < 5\,\degree$, and the $\alpha$-enhancement increases at higher latitudes.

5.3. Kinematics as a Function of [Fe/H] into the Midplane

We now examine the APOGEE bulge stars as a function of [Fe/H], by dividing the stars into the same four metallicity bins used in the ARGOS analysis of the bulge (Ness et al. 2013b): A ([Fe/H] $> 0$), B ($-0.5 \leq$ [Fe/H] $< 0$), C ($-1.0 \leq$ [Fe/H] $< -0.5$), and D ([Fe/H] $\leq -1.0$). We choose these relatively coarse bins for this analysis because, while the spatial coverage of the fields is extensive in the plane, there are few stars per field. We refer to these metallicity bins as components or populations, but this is not intended to tie them to a separate formation origin. Rather, this is intended as a tool to understand the changing kinematics of stars as a function of [Fe/H].

Figure 16 presents the rotation and dispersion profiles for stars out to $l = 35\,\degree$, for all latitudes combined and across longitude bins of $\pm 3\,\degree$, for each of these four metallicity bins. The rotation is shown in the top panel and the dispersion in the bottom panel. These stars show different kinematics as a function of [Fe/H]. Stars in components A and B, with [Fe/H] $> -0.5$, show similar rotation and dispersion, though the most metal-rich stars are the kinematically coolest population. Stars in metallicity bin C, with $-1.0 < \text{[Fe/H]} < -0.5$, show a similar, but slower, rotation profile to the more metal-rich stars and a hotter and flatter dispersion profile than stars with [Fe/H] $> -0.5$ as a function of longitude. All stars with [Fe/H] $> -1.0$ (in components A, B, C) show a rotation and dispersion that transition smoothly from the boxy bulge region inside of $l < 15\,\degree$ to the disk.

Stars with [Fe/H] $> -0.5$ have a dispersion and rotation profile that is seen in the N-body models of boxy/peanut bulges (shown in Figure 8). Stars with [Fe/H] $< -0.5$ have a dissimilar dispersion profile to the N-body models. However, the rotation and dispersion profile of stars with $-1.0 < \text{[Fe/H]} < -0.5$ and their smooth transition into the disk suggest that these stars belong to the thicker part of the disk in the inner region that is not part of the boxy/peanut morphology (e.g., Ness et al. 2013b; Di Matteo et al. 2015).

Stars in metallicity bin D, with [Fe/H] $< -1.0$, show the slowest rotation. The rotation and dispersion of these metal-poor stars diverge most significantly from those of the more metal-rich stars outside of the bulge at $l > 10\,\degree$. Notably, outside of the boxy bulge region this population of stars with [Fe/H] $< -1.0$ does not appear to be rotating, whereas inside of $l < 15\,\degree$ the stars with [Fe/H] are slowly rotating, on the order of about 50% of the more metal-rich stars with [Fe/H] $> -0.5$.

The consistency in the rotation of stars with [Fe/H] $> -1.0$ and smooth transition out in longitude suggest that these stars have a common origin, from the disk, though the more metal-poor stars with [Fe/H] $< -0.5$ do not have a dispersion profile that matches N-body models of a boxy bulge formed via instabilities. Stars with [Fe/H] $< -1.0$ are likely halo stars in the inner region or perhaps a small population of stars in the bulge that is unique to the bulge (formed via mergers or dissipational collapse). Note that stars with [Fe/H] $< -1.0$ represent only 5% of stars in the bulge, so this corresponds to a very small population.

Figures 17 and 18 present the rotation and dispersion profiles of stars in these four metallicity bins in more detail, across ($l, b$), where there are 2700, 2700, 1700, and 450 stars in each of these maps, respectively.

The APOGEE mean rotation trends for the stars in the plane ($|b| < 2\,\degree$) shown in Figures 17 and 18 are similar to the higher-latitude observations from ARGOS (Ness et al. 2013b). As seen in Figure 17, the rotation is cylindrical for populations A, B, and C and increases smoothly with longitude. Although stars in bin C show a similarly cylindrical rotation profile to stars with [Fe/H] $> -0.5$, they have an overall slower rotation profile than the more metal-rich stars in A and B. The stars in the most metal-poor bin, D, are no longer cylindrically rotating.
but display a relatively flat rotation profile with longitude across \((l, b)\).

Similarly to the bottom panel of Figure 16, the velocity dispersion maps of Figure 18 demonstrate that the stars in the metallicity bins A and B share a similar dispersion profile across \((l, b)\). The dispersion decreases with longitude and flattens for populations A–C. Inside of \(l < 10^\circ\), the dispersion decreases as a function of latitude, but is fairly constant with latitude outside of the boxy extent of the bulge. The most metal-poor stars (population D) display both high and low dispersions, though there are only a few stars in many of these fields.

6. DISCUSSION

The kinematics of the bulge mapped by APOGEE demonstrates that at latitudes \(|b| < 5^\circ\), stars with \([\text{Fe/H}] > -1.0\) have cylindrical rotation, similar to patterns farther from the plane mapped by other surveys (e.g., ARGOS, BRAVA, and GIBS), with a smooth transition in longitude out to the disk. The trends in the rotation as a function of metallicity at \(|b| < 5^\circ\) confirm and extend the findings from the ARGOS survey (determined for \(|b| > 5^\circ\)). The most metal-rich stars, with \([\text{Fe/H}] > -0.5\) (populations A and B in Section 5.3), show the fastest rotation, and the stars with
metallicities $-1.0 < [\text{Fe/H}] < -0.5$ (population C in Section 5.3), though also rotating cylindrically, have the slowest mean rotation among all stars with $[\text{Fe/H}] > -1.0$. The similar chemodynamical behavior of stars in the midplane compared to those at higher latitudes is consistent with bulge formation via instabilities of the Galactic disk and agrees well with simple $N$-body models of bulge formation from a bar that is formed in the disk.

In the kinematically hottest region, inside the boxy part of the bulge with $(|l|, |b|) < (5^\circ, 5^\circ)$, there appears to be no metallicity gradient in ($l$, $b$), as previously reported from smaller samples. This result is in contrast with higher latitudes, where the metallicity decreases by about $-0.45$ dex kpc$^{-1}$ along the minor axis (e.g., Zoccali et al. 2008; Ness et al. 2013a). A similar gradient is seen in latitude across the bulge at larger longitudes, and a weaker gradient is present across longitude, for $(|l|, |b|) > (5^\circ, 5^\circ)$ (Gonzalez et al. 2011). These gradients likely reflect the details of the redistribution of stars into the boxy/peanut bulge from the bulge-forming dynamical instabilities (e.g., Di Matteo et al. 2015) seen in $N$-body models. In these models, stars that are found in the inner bulge region typically originate from the innermost region of the disk, and this kinematically hot inner region today likely reflects a well-mixed population built from stars of a limited radial extent.

There are metal-rich, alpha-poor stars concentrated to within $b < 2^\circ$. These stars are likely part of the bar in the plane beyond the boxy extent of the bulge $l > 10^\circ$ (Wegg et al. 2015); these
stars show a smooth chemodynamical transition into the disk. The super-thin bar reported by Wegg et al. (2015) does not have any clear corresponding signature in the chemistry of APOGEE stars that differentiates them from all stars within $|b| < 2\degree$.

The chemodynamics of stars with $\mathrm{[Fe/H]} > -1.0$ suggests a common origin of formation in the disk, although stars with $-1.0 > \mathrm{[Fe/H]} > -0.5$ show kinematics that indicate that these stars are not part of the boxy structure. The small fraction ($\approx 5\%$) of metal-poor bulge stars with $\mathrm{[Fe/H]} < -1.0$ show a slower rotation and flatter dispersion that aligns them with a halo population in the innermost region or else a unique metal-poor bulge population formed at early times (see Howes et al. 2015; Koch et al. 2015).

The stars that are found at larger $(|l|, |b|)$ in the bulge, in the corners of the boxy/bulge structure, for example, at $(12\degree, 12\degree)$, are typically redistributed there from further out in the disk, and it has been demonstrated that during bulge formation the stars are preferentially redistributed according to their initial phase space (e.g., Martínez-Valpuesta & Gerhard 2013; Di Matteo et al. 2015). This redistribution effectively maps the metallicity gradient from the early disk into the bulge. The signatures of this mapping are contained in the relationship between the kinematics and $\mathrm{[Fe/H]}$ of stars as a function of $(R, z)$ today. Therefore, chemodynamical models, or models containing star formation that can track stellar ages, may provide important insights to the interpretation of the trends seen in Figures 17 and 18 and discussed in Section 5.3.

7. CONCLUSIONS

With the low-latitude data from APOGEE covering regions of the bulge never before explored with significant stellar samples, we have mapped the kinematic profile of the bulge in rotation and dispersion out to the disk. We have demonstrated that these profiles are consistent with the chemodynamical findings reported by other surveys at higher latitudes and report new details on the bulge chemodynamics. We report a population of metal-rich, alpha-poor stars that populate the plane region within $b < 2\degree$, with metal-rich stars being present at the end of the thin bar, which reaches beyond the boxy profile of the bulge, to $l \approx 25–30\degree$. These metal-rich stars at the end of the bar transition smoothly out into the disk in rotation and dispersion.

The kinematic profile of the APOGEE stars shows a dependence on $\mathrm{[Fe/H]}$ that is very similar to the trends seen in ARGOS, with stars with $\mathrm{[Fe/H]} > -1.0$ having cylindrical rotation and stars with $\mathrm{[Fe/H]} > -0.5$ showing a dispersion profile characteristic of $N$-body models of boxy bulges formed from the disk. All stars with $\mathrm{[Fe/H]} > -1.0$ show a smooth kinematic transition from the bulge region into the disk and likely have an origin in the disk. However, stars with $-1.0 < \mathrm{[Fe/H]} < -0.5$ do not appear to belong to the boxy bulge morphology given that they do not show the split red clump feature, which is the signature of the boxy/peanut structure (Ness et al. 2012; Uttenthaler et al. 2012). Furthermore, the dispersion profile of stars with $\mathrm{[Fe/H]} < -0.5$ does not match the $N$-body models of boxy peanut bulges. The small fraction of stars (5%) with $\mathrm{[Fe/H]} < -1.0$ show kinematics...
that are dissimilar in both rotation and dispersion to $N$-body models and the more metal-rich stars. These metal-poor stars do not transition smoothly in kinematics out into the disk.

Our kinematic maps for the bulge show remarkable similarity to those of other galaxies with boxy/peanut bulges. This result is perhaps not surprising, given the Milky Way’s boxy/peanut bulge profile that has been revealed in detail in the recent work of Wegg et al. (2015) and Portail et al. (2015) as a typical boxy and off-center X-shaped structure. The larger data set that will be obtained with APOGEE-2 will enable a more detailed and further examination of the chemo-kinematics of the bulge. This includes mapping trends in $[\text{Fe} / \text{H}]$ and also individual elemental abundances and also testing the symmetry around the major axis and at the ends of the bar, at both positive and negative longitudes.

We thank the referee for a constructive report that has improved the clarity of this work.

We thank Ortwin Gerhard (MPE), Chris Wegg (MPE), Matthieu Portail (MPE), Eddie Schlafly (MPIA), and Andy Casey (IoA, Cambridge) for helpful discussions. The research has received funding from the European Research Council under the European Union’s Seventh Framework Programme (FP 7) ERC Grant Agreement no. [321035].

E.A. acknowledges financial support from the CNES (Centre National d’Études Spatiales—France) and from the People Programme (Marie Curie Actions) of the European Union’s Seventh Framework Programme FP7/2007–2013/ to the DAGAL network under REA grant agreement number PITN-GA-2011–289313.

G.Z. was supported by an NSF Astronomy & Astrophysics Postdoctoral Fellowship under Award No. AST-1203017.

J.A.J. acknowledges support from NSF AST-1211853.

Funding for SDSS-III has been provided by the Alfred P. Sloan Foundation, the Participating Institutions, the National Science Foundation, and the U.S. Department of Energy Office of Science. The SDSS-III Web site is http://www.sdss3.org/.

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