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The Malliavin derivative and compactness: an application to a degenerate PDE-SDE coupling

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Abstract

Compactness is one of the most versatile tools in the analysis of nonlinear PDEs and systems. Usually, compactness is established by means of some embedding theorem between functional spaces. Such theorems, in turn, rely on appropriate estimates for a function and its derivatives. While a similar result based on simultaneous estimates for the Malliavin and weak Sobolev derivatives is available for the Wiener-Sobolev spaces, it seems that it has not yet been widely used in the analysis of highly nonlinear parabolic problems with stochasticity. In the present work we apply this result in order to study compactness, existence of global solutions, and, as a by-product, the convergence of a semi-discretisation scheme for a prototypical degenerate PDE-SDE coupling.

Keywords: compactness; degenerate diffusion; existence of solutions; Malliavin derivative; parabolic system; PDE-SDE coupling; weak-strong solution.
MSC2010: 34F05, 35B45, 35B65, 35K59, 35K65, 60H07, 60H10.

1 Introduction

Compactness is one of the most versatile tools in the analysis of nonlinear equations and systems. Usually, it is established by means of a compactness criterion for a particular functional space. Well-known examples include: the Rellich-Kondrachov theorem, the Lions-Aubin lemma, the Arzelà-Ascoli and Fréchet-Kolmogorov theorems. Such results rely on appropriate estimates of a function and its classical or weak derivatives or, more generally, of its increments. These theorems are most helpful instruments in the study of deterministic differential equations and systems in Hölder and Sobolev spaces. A closely related result was established for Wiener-Sobolev spaces by Bally and Saussereau [2]. It is based on simultaneous estimates for the Malliavin and weak Sobolev derivatives. It seems, however, that this criterion has not yet been widely used in the analysis of highly nonlinear parabolic problems with stochasticity. Indeed, standard approaches mostly rely on some monotonicity of the elliptic part (see [10, 13, 17] and references therein) which often fails to hold for strongly coupled systems. The approach based on a priori estimates and a compact embedding has several advantages. For instance, it allows to treat rather general classes of complicated problems by approximating them with better-studied, more regular ones, following the so-called compactness method [12]. In the present work we adopt this scheme. Namely, we apply the result by Bally and Saussereau in order to study compactness and existence of global solutions for a prototypical degenerate PDE-SDE coupling. The proof of existence is based on the semi-discretisation method. As a by-product, it justifies the convergence of a semi-discretisation scheme for our problem.

This paper is organised as follows. First, we introduce our model system in Section 2, fix some notations in Section 3, and state the main results in Section 4. We then establish in Sections 5-6 a set of uniform a priori estimates for the solutions of our system. While estimates in Subsections 5.1-5.4 and Section 6 are rather standard, those in Subsections 5.5-5.6 are new and more involved. The a priori estimates lie at the core of the proof of compactness in Section 7. In Section 8 we introduce a spatial discretisation scheme and study its convergence in the nondegenerate case. Finally, we use compactness in order to prove the existence of global solutions to the original degenerate problem in Section 9.

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2 The model

Let $T$ be a positive number and $O$ be a smooth bounded domain in $\mathbb{R}^N$. Also, let $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \in [0,T]}, \mathbb{P})$ be a filtered probability space on which is defined a Wiener process $(W(t))_{t \in [0,T]}$. We assume that $(\mathcal{F}_t)_{t \in [0,T]}$ is the usual completion of the natural filtration of $(W(t))_{t \in [0,T]}$. In this setting, we consider the following system of a random porous medium equation (random PME, RPME)

\[
\begin{aligned}
\begin{cases}
\partial_t \beta(c) = \Delta c + f(c, y) & \text{in } \Omega, \\
\partial_n c = 0 & \text{in } \partial \Omega, \\
\end{cases}
\end{aligned}
\]

and an Itô SDE

\[
\begin{aligned}
\begin{cases}
dy = a(y) dW + b(c, y) dt & \text{in } (0, T] \times \Omega, \\
y = y_0 & \text{in } \{0\} \times \Omega.
\end{cases}
\end{aligned}
\]

Couplings of a PDE, however with a linear diffusion, with an SDE have lately emerged in the multiscale tumor modelling [9]. In that work, the PDE- and SDE-variables represented, respectively, the intracellular and extracellular proton dynamics in a tumor. In the present case, the variables $c$ and $y$ could be seen, e.g., as the tumor density and concentration of the intracellular protons, respectively (see also [6, 7] for models based on PDE-RODE-couplings). Thereby, the (one-dimensional) Wiener process in SDE (2.2a) captures some stochastic fluctuations in the intracellular proton dynamics (see [9] and references therein). Through the coupling terms $f$ and $b$, these fluctuations influence indirectly the dynamics of the cancer cells on the macro scale as well.

The class of functions $\beta$ considered in the present study (see Assumptions 4.1 below) includes as a particular case

$$\beta(c) = c^{\frac{1}{m}}$$

for some $m > 1$.

For such $\beta$, we can transform the macroscopic PDE (2.1a) into the equation

$$\partial_t u = \Delta u^m + f(u^m, y),$$

where

$$u := c^{\frac{1}{m}}.$$

Thus, switching to the new variable $u$, we regain a RPME with a source term, written in the conventional form. Porous medium equations are standard examples of degenerate-diffusion equations. As in our recently proposed deterministic models, [20, 21], degeneracy accounts for a finite speed of propagation of a tumor. While in [20, 21] we assumed the diffusion coefficient to be degenerate not only in $c$, but also in another variable, here we consider a simpler case, with the diffusion being of the porous medium type.

3 Basic notation and functional spaces

We denote $\mathbb{R}^+ := (0, \infty)$, $\mathbb{R}^+_0 := [0, \infty)$.

For a Lebesgue measurable set $E$ we denote by $|E|$ its Lebesgue measure. The space dimension depends on the context. The integral average of an integrable function $f : E \to \mathbb{R}$ is defined via

$$\int_E f(x) \, dx := \frac{1}{|E|} \int_E f(x) \, dx.$$

We assume a smooth bounded domain $O \subset \mathbb{R}^N$, $N \in \mathbb{N}$, to be given. The outward unit normal vector on the boundary of $O$ we denote by $\nu$.

The derivative of a function $u$ of one real variable is denoted by $u'$. Partial derivatives in the classical or distributional sense with respect to a variable $z$ are denoted by $\partial_z$. The variable $z$ can, for example, be the 'time' variable $t \in \mathbb{R}^+_0$ or a component of the 'spatial' variable $x \in \bar{O}$. Further, $\nabla$ and $\Delta$ stand for the spatial gradient and Laplace operator, respectively.

We assume the reader to be familiar with the standard $L^p$, Sobolev, and Hölder spaces and their usual properties, as well as with the more general $L^p$ spaces of functions with values in general Banach
spaces, and with anisotropic spaces. In particular, for relatively open w.r.t. \((0, T] \times O\) (not necessarily cylindrical) sets \(Q\), we need the spaces

\[
W^{(1,2), q}(Q) := \{ u = u(t,x) \in L^q(Q) : \partial_t u, \partial_x u, \partial_{x_i} u, \partial_{x_j} u \in L^q(Q) \text{ for } i,j = 1, \ldots, N \}
\]

with a norm defined via

\[
\|u\|_{W^{(1,2), q}(Q)} := \left( \|u\|^q_{L^q(Q)} + \|\partial_t u\|^q_{L^q(Q)} + \sum_{i=1}^N \|\partial_{x_i} u\|^q_{L^q(Q)} + \sum_{i,j=1}^N \|\partial_{x_i} \partial_{x_j} u\|^q_{L^q(Q)} \right)^{\frac{1}{q}}.
\]

We recall that the Hölder coefficient for a Hölder exponent \(\gamma \in (0, 1)\) and a real-valued function \(w\) defined in a set \(A \subset \mathbb{R}^k, k \in \mathbb{N}\), is given by

\[
|w|_{C^\gamma(A)} := \sup_{x,y \in A, x \neq y} \frac{|w(x) - w(y)|}{|x - y|^{\gamma}}.
\]

By \(\partial A\) we denote the topological boundary of a set \(A \subset \mathbb{R}^k, k \in \mathbb{N}\). For a set \(Q \subset [0, T] \times \overline{O}\), we call

\[
\partial Q \setminus \{(T) \times O\}
\]

the parabolic boundary of \(Q\).

Further, let a filtered probability space \((\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbb{P})\) on which is defined a Wiener process \((W(t))_{t \geq 0}\) be given. We assume that \((\mathcal{F}_t)_{t \geq 0}\) is the usual completion of the natural filtration of \((W(t))_{t \geq 0}\). The corresponding Itô differential is denoted by \(dW\). We presuppose that the reader is familiar with some Itô and Malliavin calculi. In particular, we assume such standard results as: the Itô isometry (see, e.g., [14, Chapter 1 Theorem 7.1]) for \(p = 2\), a version of the Burkholder-Gundy-Davis inequality as stated in [14, Chapter 1 Theorem 7.1], as well as the Kolmogorov’s continuity criterion (see, e.g., [4, Chapter 1 Theorem 3.1] and the subsequent remark) to be known. For \(T > 0\), the Malliavin derivative of an \(\mathcal{F}_T\)-measurable \(c\) is denoted by \((D_c)_{c \in [0,T]}\). We make use of the following properties of the Malliavin derivative: the chain rule [16, Proposition 1.2.3] and a result on the weak differentiability of solutions to Itô SDEs (see, e.g., [16, Theorem 2.2.1] and the subsequent observation). We refer to [4, 14, 16] for more details on the calculus for stochastic processes.

To ease the notation while dealing with purely PDE (SDE) properties which hold \(\mathbb{P}\)-a.s. in \(\Omega\) (a.e. in \(O\)), we sometimes drop the dependence upon variable \(\omega\) (variable \(x\)) and write, for example, \(c\) instead of \(c(\cdot, \omega, \cdot)\) \((c(\cdot, \cdot, x))\). Moreover, for a stochastic process \(u : [0, T] \times \Omega \to V\), where \(V\) is a space of functions defined for \(x \in \overline{D}\), we often write \(u(t)\) instead of \(u(t,\cdot)\).

Finally, we make the following two useful conventions. Firstly, for all indices \(i, C_i\) or \(\alpha_i\) denotes a non-negative constant or, alternatively, a non-negative function, which is non-decreasing in each of its arguments. Secondly, we assume that the reappearing numbers \(T, r, t, h_1\), and \(h_2\) always satisfy

\[
0 < r < t \leq T, \quad 0 < h_1 \leq T - t, \quad 0 < h_2 \leq t - r.
\]

4 Problem setting and main result

We make the following assumptions on the problem parameters.

Assumptions 4.1.

1. \(O\) is a smooth bounded domain in \(\mathbb{R}^N, N \in \mathbb{N}\).

2. Function \(\beta : \mathbb{R}_0^+ \to \mathbb{R}_0^+\) satisfies for some constants \(m_1 > 1, m_2 \geq 1,\) and \(M, \mu > 0\) the conditions

\[
\beta\text{ has an inverse function }\beta^{(-1)}, \quad \beta(0) = 0, \quad (4.1)
\]

\[
\beta \in C^1(\mathbb{R}^+) \cap C(\mathbb{R}_0^+), \quad \beta' > 0 \text{ in } \mathbb{R}^+, \quad (4.2)
\]

\[
\beta' \text{ is decreasing in } \mathbb{R}^+, \quad (4.3)
\]

\[
(\beta')^{-1} \in C^{1, \frac{1}{\max(\mu, M)}}(\mathbb{R}_0^+), \quad (4.4)
\]

\[
c^{1, \frac{1}{\max(\mu, M)}} \beta'(c) \leq \mu \quad \text{for } c \in (0, M]. \quad (4.5)
\]
3. Functions \( f, b : \mathbb{R}_0^+ \times \mathbb{R}_0^+ \to \mathbb{R} \), and \( a : \mathbb{R}_0^+ \to \mathbb{R} \) satisfy the conditions
\[
\sup_{c, y \in \mathbb{R}_0^+} \frac{f(c, y)}{\beta(c) + 1} < \infty, \tag{4.6}
\]
\[
f \in C^1(\mathbb{R}_0^+ \times \mathbb{R}_0^+), \quad \hat{c}_f, \hat{c}_y f \in C(\mathbb{R}_0^+; C_b(\mathbb{R}_0^+)), \tag{4.7}
\]
\[
a \in C^1(\mathbb{R}_0^+), \quad a' \in C_b(\mathbb{R}_0^+), \tag{4.8}
\]
\[
b \in C^1(\mathbb{R}_0^+ \times \mathbb{R}_0^+), \quad \hat{c}_b, \hat{c}_y b \in C(\mathbb{R}_0^+; C_b(\mathbb{R}_0^+)). \tag{4.9}
\]

Moreover, \( f, a, \) and \( b \) also satisfy
\[
f(0, y) \geq 0 \text{ for all } y \in \mathbb{R}_0^+, \quad a(0) = 0, \quad b(c, 0) \geq 0 \text{ for all } c \in \mathbb{R}_0^+. \tag{4.10}
\]

4. The initial data \( c_0, y_0 : \Omega \to \mathbb{R}_0^+ \) satisfy
\[
c_0 \in L^2(\Omega) \cap H^1(\Omega), \tag{4.11}
\]
\[
y_0 \in H^1(\Omega), \tag{4.12}
\]
as well as
\[
c_0, y_0 \geq 0 \quad \text{a.e. in } O. \tag{4.13}
\]

**Definition 4.2** (Weak-strong solution). Let Assumptions 4.1 be satisfied. Let \( T > 0 \). We call a pair \((c, y) : [0, T] \times \Omega \to \mathbb{R}_0^+ \times \mathbb{R}_0^+\) a weak-strong local solution of system (2.1)-(2.2) if it holds that:

1. \((c, y) : [0, T] \times \Omega \to L^2(\Omega)\) is an adapted process;
2. \( c \in L^2(\Omega; L^2([0, T]; L^2(\Omega))) \), \( \hat{c} c \in L^2(\Omega; L^2([0, T]; L^2(\Omega))) \);
3. \( c \in L^2(\Omega; L^2([0, T]; H^1(\Omega))) \), \( \hat{c} c \in L^2(\Omega; L^2([0, T]; H^1(\Omega))) \);
4. \( D_r c \in L^2((r, T); L^2(\Omega; L^2(\Omega))) \), \( \hat{c} D_r \beta(c) \in L^2((r, T) \times \Omega; H^{-2}(\Omega)) \) for all \( r \in [0, T] \);
5. \( y \in L^2([0, T]; L^2(\Omega; H^1(\Omega))) \cap C^{1/2}([0, T]; L^q(\Omega; L^q(\Omega))) \) for some \( q > 2 \);
6. \( D_r y \in L^2((r, T); L^2(\Omega; L^2(\Omega))) \) for all \( r \in [0, T] \);
7. \((c, y)\) satisfies
\[
\begin{aligned}
- \beta(c(\cdot), \cdot) v_{L^2(\Omega)}(t) & \xi(0) = \int_0^T \langle \beta(c(s), v) \rangle_{L^2(\Omega)} \frac{d\xi}{ds}(s) ds \\
& = - \int_0^T \langle \nabla c(s), \nabla v \rangle_{L^2(\Omega)}(s) \xi(s) ds + \langle f(c, y)(s), v \rangle_{L^2(\Omega)}(s) ds
\end{aligned}
\]
for all \( v \in H^1_0(\Omega) \) \( (v \in H^1(\Omega)) \) for all \( \xi \in C^1_{0}[0, T] \) a.s. in \( \Omega \), \( \tag{4.14a} \)
\[
y(t) = y_0 + \int_0^t a(y(s)) dW(s) + \int_0^t b(c, y(s)) ds \text{ in } L^2(\Omega) \text{ for all } t \in [0, T] \text{ a.s. in } \Omega. \tag{4.14b}
\]

If \((c, y)\) is a local solution for all \( T > 0 \), then we call it a global solution.

**Remark 4.3** (Weak-strong solution). We call a solution \((c, y)\) from Definition 4.2 a weak-strong solution since it satisfies the PDE (2.1) in a weak PDE-sense and the SDE (2.2) in the strong SDE-sense.

**Remark 4.4** (Continuity of sample paths).

1. Conditions \( c \in L^2(\Omega; L^2([0, T]; H^1(\Omega))) \) and \( \hat{c} c \in L^2(\Omega; L^2([0, T]; L^2(\Omega))) \) imply that \( c(\cdot, \cdot, \omega) \in H^1([0, T]; L^2(\Omega)) \) a.s. in \( \Omega \). Standard result [12, Chapter 1 Lemma 1.2] yields that \( c(\cdot, \cdot, \omega) \in C([0, T]; L^2(\Omega)) \) a.s. in \( \Omega \).

2. Since \( q > 2 \), condition \( y \in C^{1/2}([0, T]; L^q(\Omega; L^q(\Omega))) \) implies that \( y(\cdot, \cdot, \cdot) \in C([0, T]; L^q(\Omega)) \) a.s. in \( \Omega \). This is a direct consequence of the Kolmogorov’s continuity criterion (see [3, Chapter 1 Problem 3.1] and the subsequent remark). Below we choose \( q \) as in the Sobolev embedding theorem, i.e., such that \( H^1(\Omega) \subset L^q(\Omega) \).
For arbitrary $T > 0$ and $R = (R_0, R_1), R_0, R_1 > 0$, we introduce the parameter set

$$\mathcal{P}(T, R) := \{ (\beta, f, a, b, c_0, y_0) \mid \text{as in Assumptions 4.1 and such that } \|c_0\|_{L^p(O)}, \sup_{c, y \in \mathbb{R}^k} \frac{f(c, y)}{\beta(c) + 1} \leq R_0, \sup_{c \in [0, R_2(T, R_0)], \|c\|_{C^1([0, R_2(T, R_0)])}} c^{1 - \frac{1}{p}} \beta'(c), \|c \partial_c f, \partial_y f, \partial_c b, \partial_y b\|_{C^1([0, R_2(T, R_0)] \times \mathbb{R}^k)} \leq R_1, \|a'\|_{C^1([0, R_0])}, b(0, 0), \|c_0, y_0\|_{H^1(O)} \leq R_1 \}$$

where

$$R_2(T, R_0) := \beta^{(-1)}\left( e^{Tr_0} (\beta(R_0) + 1) - 1 \right) \tag{4.15}$$

is an upper bound for solution component $c$ if the parameters belong to $\mathcal{P}(T, R)$ (see estimate (5.4) below). We define the corresponding sets of solutions and of their components via

$$\mathcal{U}(T, R) := \{ (c, y) \mid \text{is a weak-strong solution corresponding to } \beta, f, a, b, c_0, y_0 \in \mathcal{P}(T, R) \},$$

$$\mathcal{C}(T, R) := \{ (c) \mid (c, y) \in \mathcal{U}(T, R) \},$$

$$\mathcal{Y}(T, R) := \{ (y) \mid (c, y) \in \mathcal{U}(T, R) \}.$$ 

Now we are ready to formulate our compactness result:

**Theorem 4.5 (Compactness).** Let Assumptions 4.1 1.-3. be satisfied. Then for all $T > 0$ and $R = (R_0, R_1), R_0, R_1 > 0$, it holds that

$$\mathcal{C}(T, R) \text{ and } \mathcal{Y}(T, R) \text{ are precompact in } L^2((0, T) \times O \times \Omega).$$

We prove this theorem in Section 7 and then use it in Section 9 in order to establish the following result which deals with the existence of solutions:

**Theorem 4.6 (Existence of a weak-strong solution).** Let Assumptions 4.1 be satisfied. Then there exists a weak-strong global solution in terms of Definition 4.2 to system (2.1)-(2.2).

**Remark 4.7 (Uniqueness).** The uniqueness of solutions to (2.1)-(2.2) holds as well. It can be proved in a standard way by exploiting the monotonicity of $\beta$ without requiring the solutions to be differentiable in the Malliavin sense.

**Remark 4.8 (Notation).** We make the following useful convention: the statement that a quantity (a constant or a function) depends on the parameters of the problem means that it depends upon the space dimension $N$, domain $O$, constants $T, m_1, m_2, \mu$, and $q$, and the structure of the initial values $c_0$ and $y_0$ and of the coefficient functions $\beta, f, a$ and $b$ (the latter means their norms etc. which appear in the definition of the parameter set $\mathcal{P}(T, R)$).

Moreover, dependence upon these parameters is mostly not indicated in an explicit way.

## 5 A priori estimates for the RPME (2.1)

### 5.1 Standard PDE estimates for $c$

Equation (2.1a) is a.s. in $\Omega$ a standard PDE in $(0, T] \times O$. This allows us to derive in a standard way several basic estimates for $c$ which hold irrespectively of $\omega$. To begin with, we multiply (2.1a) by $p(\beta(c) + 1)^{p-1}$ for an arbitrary $p > 1$ and integrate by parts over $O$ using the boundary conditions. We thus obtain with the help of assumptions (4.2) and (4.6) that

$$\frac{d}{dt}\|\beta(c) + 1\|_{L^p(O)} = - p(p - 1) \int_O c \beta'(c)(\beta(c) + 1)^{p-2} |\nabla c|^2 \, dx + p \int_O f(c, y)(\beta(c) + 1)^{p-1} \, dx$$

$$\leq pC_1 \|\beta(c) + 1\|_{L^p(O)},$$

where

$$C_1 := \sup_{c, y \in \mathbb{R}^k} \frac{f(c, y)}{\beta(c) + 1}.$$
Applying the Gronwall lemma to (5.1) and taking the $\frac{1}{p}$ power on both sides of the resulting inequality, we obtain that
\[ \|\beta(c)(t) + 1\|_{L^p(O)} \leq e^{tC_1} \|\beta(c_0) + 1\|_{L^p(O)}. \] (5.2)
In the limit as $p \to \infty$, estimate (5.2) and the fact that $\beta$ is increasing yield that
\[ \|\beta(c)(t) + 1\|_{L^\infty(O)} \leq e^{tC_1} \|\beta(c_0) + 1\|_{L^\infty(O)} \]
\[ \leq e^{tC_1} \left( \|c_0\|_{L^\infty(O)} + 1 \right). \] (5.3)
Consequently, we arrive with (5.3) and assumption (4.1) at the estimate
\[ \|c\|_{L^\infty((0,T) \times O \times \Omega)} \leq \beta^{-1} \left( e^{tC_1} \left( \|c_0\|_{L^\infty(O)} + 1 \right) - 1 \right), \] (5.4)
where
\[ R_0 := \max\{C_1, \|c_0\|_{L^\infty(O)}\}, \]
and $R_2$ was defined in (4.15).

**Remark 5.1.** Due to estimate (5.4), it suffices to consider the coefficient functions for $c \in [0, R_2(T, R_0)]$ only.

Another standard estimate for $c$ as solution to the PME (2.1) is obtained by multiplying by $\hat{\gamma}_t c$ and integrating by parts over $O$ using the boundary conditions and then over $[0, t]$ for $t \in (0, T]$. This implies due to (5.4), assumptions (4.3) and (4.11), and the Young inequality that
\[ C_2 \int_0^t \|\hat{\gamma}_t c(s)\|_{L^2(O)}^2 \, ds \leq \beta' \left( \|c\|_{L^\infty((0,T) \times O \times \Omega)} \right) \int_0^t \|\hat{\gamma}_t c(s)\|_{L^2(O)}^2 \, ds \]
\[ \leq \int_0^t \int_0^t \beta'(x) |\hat{\gamma}_t c|^2(s) \, dx \, ds \]
\[ = - \|\nabla c(t)\|_{L^2(O)}^2 + \|\nabla c_0\|_{L^2(O)}^2 + \int_0^t \int_0^s f(c, y) \hat{\gamma}_t c(s) \, dx \, ds \]
\[ \leq - \|\nabla c(t)\|_{L^2(O)}^2 + C_2 \int_0^t \|\hat{\gamma}_t c(s)\|_{L^2(O)}^2 \, ds + C_3. \] (5.5)
Consequently, we obtain that
\[ \int_0^t \|\hat{\gamma}_t c(s)\|_{L^2(O)}^2 \, ds, \|\nabla c(t)\|_{L^2(O)}^2 \leq C_4 \text{ for all } t \in [0, T]. \] (5.5)
Altogether, estimates (5.4) and (5.5) yield the first group of estimates for norms of $c$:
\[
\text{a.s. in } \Omega \quad \begin{cases} \\
\|c\|_{L^\infty((0,T) \times O)} \leq C_5, \\
\|\nabla c\|_{L^\infty((0,T);L^2(O))} \leq C_5, \\
|\hat{\gamma}_t c|_{L^2((0,T);L^2(O))} \leq C_5.
\end{cases} \] (5.6)
(5.7)
(5.8)
In particular, assumption (4.6) together with estimate (5.6) imply that
\[ \|f(c)\|_{L^\infty((0,T) \times O)} \leq C_6 \quad \text{a.s. in } \Omega. \] (5.9)

**5.2 Estimate for the Malliavin derivative $D_t c$**

The next step is to apply the Malliavin derivative operator $D_t$ on both sides of the integrated form of equation (2.1), the integro-differential equation (4.14a). Using the chain rule and the locality property $(D_{t_1}, F(t_2) = 0$ for $t_2 < t_1$ for $F$ adapted) of the operator $D_t$, we compute that
\[ \beta'(c) D_t \beta(c)(t) = D_t \beta(c)(t) \]
\[ = D_t \Delta \int_0^t c(s) \, ds + D_t \int_0^t f(c, y)(s) \, ds \]
\[ = \Delta \int_r^t D_r c(s) \, ds + \int_r^t D_r f(c, y)(s) \, ds. \tag{5.10} \]

Multiplying (5.10) by \( D_r c(t) \) and integrating by parts over \( O \) using the boundary conditions and then over \([r, t]\), we obtain using assumption (4.3), estimate (5.6), and the Young inequality that

\[
\begin{align*}
C_7 \int_r^t \|D_r c(s)\|^2_{L^2(O)} \, ds \\
\leq \int_r^t \int_O \beta'(c) |D_r c|^2(s) \, dx \, ds \\
= -\frac{1}{2} \| \nabla \int_r^t D_r c(s) \|_{L^2(O)}^2 \| + \int_r^t \int_O D_r c(s) \int_r^t D_r f(c, y)(\tau) \, d\tau \, dx \, ds \\
\leq \frac{C_7}{2} \int_r^t \|D_r c(s)\|^2_{L^2(O)} \, ds + C_8 \int_r^t \int_r^t \|D_r f(c, y)(\tau)\|^2_{L^2(O)} \, d\tau \, ds.
\end{align*}
\]

Thus, we arrive at the inequality

\[ \int_r^t \|D_r c(s)\|^2_{L^2(O)} \, ds \leq C_9 \int_r^t \int_r^t \|D_r f(c, y)(\tau)\|^2_{L^2(O)} \, d\tau \, ds. \tag{5.11} \]

In order to estimate the term on the right hand side of (5.11), we use the chain rule and assumption (4.7). We thus obtain that

\[ \|D_r f(c, y)(t)\|^2_{L^2(O)} \leq C_{10} \|D_r c(t)\|^2_{L^2(O)} + C_{10} \|D_r y(t)\|^2_{L^2(O)}. \tag{5.12} \]

Combining (5.11) and (5.12) yields that

\[ \int_r^t \|D_r c(s)\|^2_{L^2(O)} \, ds \leq C_{11} \int_r^t \int_r^t \|D_r c(\tau)\|^2_{L^2(O)} \, d\tau \, ds + C_{11} \int_r^t \int_r^t \|D_r y(\tau)\|^2_{L^2(O)} \, d\tau \, ds. \tag{5.13} \]

Application of the Gronwall lemma to (5.13) yields that

\[
\begin{align*}
\int_r^t \|D_r c(s)\|^2_{L^2(O)} \, ds &\leq C_{12} \int_r^t \int_r^t \|D_r y(\tau)\|^2_{L^2(O)} \, d\tau \, ds \\
&\leq C_{13} \int_r^t \|D_r y(s)\|^2_{L^2(O)} \, ds \quad \text{a.s. in } \Omega. \tag{5.14}
\end{align*}
\]

Moreover, estimates (5.12) and (5.14) imply that

\[ \int_r^t \|D_r f(c, y)(s)\|^2_{L^2(O)} \, ds \leq C_{14} \int_r^t \|D_r y(s)\|^2_{L^2(O)} \, ds \quad \text{a.s. in } \Omega. \tag{5.15} \]

### 5.3 Estimate for \( \partial_t D_r \beta(c) \)

Applying the differential operator \( D_r \) to the original equation (2.1a), we have that

\[ \partial_t D_r \beta(c) = \Delta D_r c + D_r f(c, y) \tag{5.16} \]

Combining estimates (5.14) and (5.15) with the Sobolev inequality, we conclude from (5.16) that

\[
\begin{align*}
\int_r^t \|\partial_t D_r \beta(c)(s)\|^2_{H^{-2}(O)} \, ds &\leq C_{15} \int_r^t \|D_r c(s)\|^2_{L^2(O)} + \|D_r f(c, y)(s)\|^2_{H^{-2}(O)} \, ds \\
&\leq C_{16} \int_r^t \|D_r c(s)\|^2_{L^2(O)} + \|D_r f(c, y)(s)\|^2_{L^2(O)} \, ds \\
&\leq C_{17} \int_r^t \|D_r y(s)\|^2_{L^2(O)} \, ds \quad \text{a.s. in } \Omega. \tag{5.17}
\end{align*}
\]
5.4 Estimate for $D_{r+h_2}c - D_r c$

Inserting $r + h_2$ in place of $r$ in (5.10), we have (recall that $D_i F(t_2) = 0$ for $t_2 < t_1$ for $F$ adapted) that

$$\beta'(c)D_{r+h_2}c(t) = \Delta \int_r^{t} D_{r+h_2}c(s) \, ds + \int_r^{t} D_{r+h_2}f(c, y)(s) \, ds,$$

$$= \Delta \int_r^{t} D_{r+h_2}c(s) \, ds + \int_r^{t} D_{r+h_2}f(c, y)(s) \, ds.$$  

Subtracting (5.18) and (5.10), we obtain that

$$\beta'(c)(D_{r+h_2}c - D_r c)(t) = \Delta \int_r^{t} (D_{r+h_2}c - D_r c)(s) \, ds + \int_r^{t} (D_{r+h_2}f(c, y) - D_r f(c, y))(s) \, ds.$$  

Arguing for (5.19) as we just did for (5.10) above, we obtain that

$$\int_r^{t} \|(D_{r+h_2}c - D_r c)(s)\|_{L^2(O)}^2 \, ds \leq C_{18} \int_r^{t} \|(D_{r+h_2}y - D_r y)(s)\|_{L^2(O)}^2 \, ds \quad \text{a.s. in } \Omega.$$  

5.5 A regularising transformation for solutions of a PME

In this sequel, we deal with some purely PDE properties of (2.1a). It is well-known that solutions of a PME like (2.1a) are, in general, only weak-strong solutions if $c \neq 0$ and is not strictly separated from zero. In particular, $\Delta c$ is generally not even $L^2$-bounded. Still, it is well-understood [1, 3, 5, 8, 22] that under reasonable assumptions on $f$ the solution is at least locally Hölder continuous. Following our idea from [19], we show how the information on the (local) Hölder continuity of a solution function can be used in order to transform this function into a smooth one by means of a smooth and strictly increasing function which depends only upon the parameters of the problem. We believe that this result is of interest by itself. We then use this transformation for the compactness proof, see Section 7. Before we begin with a construction for solutions of a PME, let us consider a simple motivating example.

**Example 5.2.** Denote by $B_r$, $r > 0$, the closed ball of radius $r$ in $\mathbb{R}^N$ centred at the origin. Let

$$w : B_1 \to [0, 1], \quad w(x) = |x|^\gamma \text{ for some } \gamma \in (0, 1).$$

It is well-known that $w \in C^2(B_1 \setminus \{0\}) \cap C^\gamma(B_1)$, but $w \notin C^2(B_1)$. Since

$$|\partial_x \partial_x^2 w(x)| \leq \gamma(\gamma - 1)|x|^{\gamma - 2} \text{ for all } x \in B_1 \setminus \{0\}$$

and

$$\{w > k\} = B_k^+ \text{ for all } k \in (0, 1],$$

we have that

$$\|\partial_x \partial_x^2 w\|_{C(\{w > k\})} \leq \gamma(\gamma - 1)k^{1-\frac{\gamma}{2}} \text{ for all } k \in (0, 1].$$  

Set

$$\varphi : [0, 1] \to \mathbb{R}^+, \quad \varphi(k) := \left(\gamma(\gamma - 1)k^{1-\frac{\gamma}{2}}\right)^{-1} = (\gamma(\gamma - 1))^{-1} k^{\frac{\gamma}{2} - 1}.$$  

We define a regularising transformation for $w$ by

$$\Phi : [0, 1] \to \mathbb{R}^+, \quad \Phi(k) := \int_0^k \int_0^{s_2} \varphi(s_1) \, ds_1 \, ds_2 = \left(\gamma(\gamma - 1) \frac{2}{\gamma} \left(\frac{2}{\gamma} + 1\right)\right)^{-1} k^{\frac{\gamma}{2} + 1}.$$  

Then, we have that

$$\Phi(w(x)) = \left(\gamma(\gamma - 1) \frac{2}{\gamma} \left(\frac{2}{\gamma} + 1\right)\right)^{-1} |x|^{2+\gamma} \text{ for all } x \in B_1.$$  

It is easy to see that $\Phi$ has the following properties:

1. $\Phi(w) \in C^2(B_1)$;
2. $\Phi$ is a strictly increasing function, so that it allows to reconstruct back $w$ from $\Phi(w)$;

3. $\Phi$ preserves the zero set of $w$ and allows to reconstruct it back;

4. $\Phi$ can be used to regularise a whole class of functions $w$ which are smooth everywhere but for their zero sets and satisfy (5.21). Thus, $\Phi$ smooths down such a function $w$ near its zero set.

Let us now apply the idea from Example 5.2 to solutions $c$ of (2.1a). In this general case, however, we cannot hope for pointwise estimates like (5.21) to hold uniformly in $\omega$. This is because the source term $f$ depends upon $y$, which is, for each $x$, a solution of an SDE. Hence, instead of using the spaces of functions which are differentiable in the classical sense, we work in anisotropic Sobolev spaces. The, possibly, irregular behaviour of $c$ at the parabolic boundary of the cylinder $(0, T) \times O$ presents yet enough difficulty. Our regularising transformation should thus be able to smooth down $c$ not only near $\{c = 0\}$, where the equation has a degeneracy, but also at the parabolic boundary

$$
\Gamma := \partial((0, T) \times O)|\{T \times O\}.
$$

For this reason, we consider $c$ on the intersections of its level sets with a decreasing family of subcylinders of $(0, T) \times O$:

$$
Q_d := \{(t, x) \in (0, T] \times O : \text{dist}((t, x), \Gamma) > d\} \quad \text{for all } d \in \left(0, \frac{1}{4} \text{diam}(O)\right).
$$

We recall that due to assumptions on $\beta$ and estimates (5.6) and (5.9) it holds (see, e.g., [8, Theorem 2.1]) that

$$
|c|_{C^{\alpha_0}(Q_d)} \leq C_{19} (d^{-1}) \quad \text{for all } d \in \left(0, \frac{1}{4} \text{diam}(O)\right) \text{ for some } \alpha_0 \in (0, 1).
$$

Thus, the Hölder constant may explode as $d \to 0$, that is, as $\Gamma$ is approached. We next divide both sides of (2.1a) by $\beta'(c)$ and thus obtain an equation in a non-divergence form:

$$
\tilde{c} \Delta c = (\beta'(c))^{-1} \Delta c + (\beta'(c))^{-1} f(c, y).
$$

Due to assumption (4.4) and estimates (5.6), (5.9), and (5.22), it holds that

$$
\|(\beta'(c))^{-1}\|_{C^{\alpha_1}(Q_d)} \leq C_{20} (d^{-1}) \quad \text{for some } \alpha_1 \in (0, 1),
$$

$$
\|(\beta'(c))^{-1} f(c, y)\|_{L^\infty(Q_d)} \leq C_{21}.
$$

In order to obtain (5.24), we used the well-known property of superpositions of Hölder continuous functions:

$$
u_1 \in C^{\gamma_1}(D), \ u_2 \in C^{\gamma_2}(u_1(D)) \quad \text{for some } \gamma_1, \gamma_2 \in (0, 1)$$

$$
\Rightarrow u_2 \circ u_1 \in C^{\gamma_1 + \gamma_2}(D) \quad \text{and } |u_2 \circ u_1|_{C^{\gamma_1 + \gamma_2}(D)} \leq |u_2|_{C^{\gamma_2}(u_1(D))} |u_1|_{C^{\gamma_1}(D)}.
$$

Let us consider for any $k \in \left(0, |c|_{L^\infty((0, T) \times O)}\right), d \in \left(0, \frac{1}{4} \text{diam}(O)\right)$ the sets $\{c > k\} \cap Q_d$ and $\{c > \frac{k}{2}\} \cap Q_d$. Using the crucial property (5.22), we deduce the following: these two sets are relatively open w.r.t. $(0, T] \times O$ and their parabolic boundaries do not intersect. Moreover, the estimate on the Hölder norm allows to estimate the distance between the parabolic boundaries from below by a positive number which depends only upon $d$, $k$, and, of course, the parameters of the problem. Now, equation (5.23) is non-degenerate in $\{c > \frac{k}{2}\}$. Therefore, we can apply standard results on local regularity for linear parabolic equations, see Theorems 9.1 and 10.1, and the remark on local estimates in Sobolev spaces at the end of §10 in [11, Chapter IV]. Considering $\{c > k\} \cap Q_d$ as a subdomain of $\{c > \frac{k}{2}\} \cap Q_d$, these regularity results can be interpreted in the following way: for each $p \in (1, \infty)$ there exists a function

$$
\varphi_p : \mathbb{R}^+ \times \left[0, \frac{1}{4} \text{diam}(O)\right] \rightarrow \mathbb{R}^+_0
$$

with the properties

1. $\varphi_p$ depends only upon $p$ and the parameters of the problem;

2. $\varphi_p(0, \cdot) = 0$, $\varphi_p(\cdot, 0) = 0$ in $\mathbb{R}_0^+$, $\varphi_p > 0$ in $\mathbb{R}^+ \times \left[0, \frac{1}{4} \text{diam}(O)\right]$;
3. \( \varphi_p \) is uniformly bounded;

4. \( \varphi_p \) is increasing in each of the two variables;

5. for each \( k \in \left( 0, \|c\|_{L^\infty((0,T) \times O)} \right) \) and \( d \in \left( 0, \frac{1}{4} \text{diam}(O) \right) \) it holds that

\[
\|c\|_{W^{1,2}(\mathbb{R}^3; (c > k) \cap Q_d)} \leq \varphi_p^{-1}(k, d);
\]  

(5.25)

6. \( \varphi_p \beta' \leq 1. \)

**Remark 5.3.** Local estimates from \([11]\) deal with cylindrical sets, which is generally not the case for a set \( \{ c > k \} \cap Q_d \). However, since the closure of such a set is a compact set, lies inside of a relatively w.r.t. \( (0,T) \times O \) open set \( \{ c > \frac{k}{2} \} \cap Q_{\frac{3}{2}}, \) and we have control on the distance between the parabolic boundaries, we can cover \( \{ c > k \} \cap Q_d \) by a finite number of sufficiently small cylinders which lie completely in \( \{ c > \frac{k}{2} \} \cap Q_{\frac{3}{2}} \) and such that their number and the distance between their parabolic boundaries and the parabolic boundary of \( \{ c > \frac{k}{2} \} \cap Q_{\frac{3}{2}} \) is bounded from below by a positive number which depends only upon \( d, k, \) and the parameters of the problem. Hence, we can apply the results from \([11]\) to each of these cylinders and subsequently sum together the resulting estimates in order to deduce (5.25) with \( \varphi \) satisfying conditions 1.-4. from above.

**Remark 5.4.** Observe that if a function \( \tilde{\varphi}_p \) satisfies conditions 1.-5. from above, we can clearly satisfy all six conditions by taking

\[
\varphi_p(k, d) := \min\{\tilde{\varphi}_p(k, d), (\beta'(k))^{-1}\}.
\]

Estimate (5.25) together with properties 1.-4. convey that \( c \) is well-behaved away from its zero set \( \{ c = 0 \} \) and the parabolic boundary of \( (0,T) \times O \), may possibly have singularities on that parabolic boundary and/or \( \{ c = 0 \} \), but, also, that we have some control on its behaviour near the singularities. Using \( \varphi_p \), we are now able to produce our regularising transformation for \( c \):

\[
\Phi_p : \mathbb{R}_0^+ \times \left[ 0, \frac{1}{4} \text{diam}(O) \right] \to \mathbb{R}_0^+,
\]

\[
\Phi_p(k, d) := \int_0^k \int_0^{s_2} \int_0^{s_2} \frac{s_1^2}{2} \varphi_p^2 \left( \frac{s_1}{2}, z_1 \right) dz_1 ds_1 ds_2.
\]

Due to properties 2. and 4. of \( \varphi_p \), we have for each \( k \in \left( 0, \|c\|_{L^\infty((0,T) \times O)} \right) \) and \( d \in \left( 0, \frac{1}{4} \text{diam}(O) \right) \) that

\[
0 \leq \partial_{k^{\alpha_1} d^{\alpha_2}} \Phi_p(k, d) \leq C_{22} \frac{k}{2} \varphi_p^2 \left( \frac{k}{2}, d \right) \text{ for } \alpha_1, \alpha_2 \in \{ 0, 1, 2 \}, \ \alpha_1 + \alpha_2 \leq 2.
\]  

(5.26)

Next, we recall that domain \( O \) has a smooth boundary. Consequently, there exists a function

\[
\gamma : [0,T] \times \overline{O} \to \left[ 0, \frac{1}{4} \text{diam}(O) \right]
\]

with the following properties:

1. \( \gamma(t,x) > 0 \text{ in } (0,T) \times O, \ \gamma(t,x) = 0 \text{ in } \Gamma; \)

2. \( \gamma \in C^2([0,T] \times \overline{O}); \)

3. there exists a number \( d_0 \in \left( 0, \frac{1}{4} \text{diam}(O) \right] \), which depends only upon the domain \( O \), such that

\[
\gamma \leq d \text{ in } \overline{Q_d} \text{ for all } d \in [0,d_0].
\]

Using the chain rule, (5.26) and the properties of \( \varphi_p \) and \( \gamma \), one readily checks that for all \( p \in (1, \infty) \) it holds that

\[
\|\Phi_p(c, \gamma)\|_{W^{1,2}((0,T) \times O)} \leq C_{23}(p),
\]

(5.27)

\[
\|\tilde{c}_t \Phi_p(c, \gamma)\beta'(c)\|_{L^P((0,T) \times O)} \leq C_{24}(p).
\]

(5.28)
Indeed, for each \(i, j \in \{1, \ldots, N\}\), \(k \in \big(0, \|c\|_{L^\infty((0,T) \times O)}\big)\) and \(d \in (0, d_0]\), it holds due to the Hölder inequality that, for instance,

\[
\left\| \partial_{x^2} \Phi_p(c, \gamma) \partial_{x^j} c \right\|_{L^p((k < c < 2k) \cap Q_d)} \leq C_{22} \left\| \frac{c}{2} \partial_{x^j} \Phi_p(c, \gamma) \right\|_{L^p((k < c < 2k) \cap Q_d)} \leq C_{22} k \varphi_p^2(k, d) \left\| \partial_{x^j} c \right\|_{L^p((k < c < 2k) \cap Q_d)} \leq C_{22} k \varphi_p^2(k, d) \left\| c \right\|_{L^p((1,2) \times \Omega)} \left\| \partial_{x^j} \right\|_{L^p((1,2) \times \Omega)} \leq C_{22} k.
\]

(5.29)

Since the constant \(C_{22}\) doesn’t depend upon \(k\), estimate (5.29) yields that

\[
\left\| \partial_{x^j} \Phi_p(c, \gamma) \partial_{x^j} c \right\|_{L^p((c > 0) \cap Q_d)} = \sum_{i=0}^\infty \left\| \partial_{x^j} \Phi_p(c, \gamma) \partial_{x^j} c \right\|_{L^p((2^{-(i+1)}|c|_{L^\infty((0,T) \times O)} < c < 2^{-i}|c|_{L^\infty((0,T) \times O)}) \cap Q_d)} \leq C_{22} \sum_{i=0}^\infty 2^{-(i+1)} = C_{22}.
\]

(5.30)

Finally, since \(\Phi_p(c, \gamma) \equiv 0\) on \(\{c = 0\}\) and \(C_{22}\) doesn’t depend upon \(d\), (5.30) implies that

\[
\left\| \partial_{x^j} \Phi_p(c, \gamma) \partial_{x^j} c \right\|_{L^p((0,T) \times O)} \leq C_{22}.
\]

Applying the chain rule to \(\Phi_p(c, \gamma)\) in order to compute the required partial derivatives and treating other resulting terms in a similar fashion leads to estimates (5.27) and (5.28).

### 5.6 Estimates for a transformation of solutions of the RPME (2.1)

Let \(p \in (1, \infty)\) and let \(\Psi_p\) be the smoothing transformation from Subsection 5.5. We now introduce yet another transformation

\[
\Psi_p : \mathbb{R}_0^+ \times \left[0, \frac{1}{4} \text{diam}(O)\right] \to \mathbb{R}_0^+,
\]

\[
\Psi_p(k, d) := \int_0^k \Phi_p(s, d)\beta'(s) \, ds.
\]

Due to assumption (4.2) and the properties of \(\Phi_p\), we have that

\[
\Psi_p \in C^1 \left( \mathbb{R}_0^+ \times \left[0, \frac{1}{4} \text{diam}(O)\right] \right), \quad 0 \leq \partial_k \Psi_p(k, d) \leq C_{25}(p, k).
\]

(5.31)

The continuity of \(\partial_k \Psi_p\) in \(\mathbb{R}_0^+ \times [0, \frac{1}{4} \text{diam}(O)]\) is a direct consequence of continuity of \(\Phi_p\) and \(\beta\). Moreover, it holds with (5.26) and the properties of \(\varphi_p\) listed above that

\[
0 \leq \Phi_p(k, d) \beta'(k) \leq \frac{k}{2} \varphi_p^2 \left( \frac{k}{2} d \right) \beta'(k) \leq C_{26}(p) k \varphi_p \beta(k) \leq C_{26}(p) k =: C_{25}(p, k).
\]

(5.32)

Estimate (5.32) yields in particular that \(\partial_k \Psi_p\) is continuous in every point of the set \(\{0\} \times \left[0, \frac{1}{4} \text{diam}(O)\right]\), as required.

Using estimates (5.6)–(5.8) and (5.31) and the chain rule, we obtain the following group of estimates:

\[
\begin{align*}
|\Psi_p(c, \gamma)|_{L^\infty((0,T) \times O)} & \leq C_{27}(p), \quad \text{a.s. in } \Omega, \quad \text{(5.33a)} \\
|\nabla_c \Psi_p(c, \gamma)|_{L^\infty((0,T) \times O)} & \leq C_{27}(p), \quad \text{(5.33b)} \\
|\partial_t \Psi_p(c, \gamma)|_{L^2((0,T) \times O)} & \leq C_{27}(p), \quad \text{(5.33c)} \\
|D_t \Psi_p(c, \gamma)|_{L^2((r,T) \times O)} & \leq C_{27}(p) \|D_t c\|_{L^2((r,T) \times O)}, \quad \text{(5.33d)} \\
|D_t + \Delta \Psi_p(c, \gamma) - D_t \Psi_p(c, \gamma)|_{L^2((r,T) \times O)} & \leq C_{27}(p) \|D_t + \Delta c - D_t c\|_{L^2((r,T) \times O)}. \quad \text{(5.33e)}
\end{align*}
\]
Next, we use the chain and product rules in order to obtain the following representation for $\partial_t D_t \Psi_p(c, \gamma)$:

$$\partial_t D_t \Psi_p(c, \gamma) = \Phi_p(c, \gamma) \partial_t \beta(c) + \partial_t \Phi_p(c, \gamma) \beta'(c) D_t c.$$  \hfill (5.34)

Let $p > \max \left\{ \frac{N}{2}, 2 \right\}$. Then, $W^2_{p}(O)$ is closed under pointwise multiplication and

$$\|uv\|_{W^2_{p}(O)} \leq C_{28}(p) \|u\|_{W^2_{p}(O)} \|v\|_{W^2_{p}(O)},$$

so that, due to the Sobolev inequality,

$$\|uv\|_{H^2_{p}(O)} \leq C_{29}(p) \|u\|_{W^2_{p}(O)} \|v\|_{W^2_{p}(O)}, \quad (5.35)$$

Using (5.35), we obtain the following estimate for the first summand on the right-hand side of (5.34):

$$\|\Phi_p(c, \gamma) \partial_t \beta(c)\|_{H^2_{p}(O)} \leq C_{30}(p) \|\Phi_p(c, \gamma)\|_{W^2_{p}(O)} \|\partial_t D_t \beta(c)\|_{H^2_{p}(O)}. \quad (5.36)$$

For the second summand, the Hölder and Sobolev inequalities together with the properties of $\Phi_p$ yield that

$$\|\partial_t \Phi_p(c, \gamma) \beta'(c) D_t c\|_{H^2_{p}(O)} \leq C_{31}(p) \|\partial_t \Phi_p(c, \gamma)\|_{L^p(r,T \times O)} \|\beta'(c)\|_{L^p(r,T \times O)} \|D_t c\|_{L^2(r,T \times O)}.$$  \hfill (5.37)

Combining (5.27), (5.28), (5.34), (5.36), and (5.37), integrating over $(r, T)$ and using the Hölder inequality, we obtain that

$$\|\partial_t D_t \Psi_p(c, \gamma)\|_{L^2(r,T \times O)} \leq C_{32}(p) \left( \|\partial_t D_t \beta(c)\|_{L^2(r,T \times O)} + \|D_t c\|_{L^2(r,T \times O)} \right) \quad \text{a.s. in } \Omega. \quad (5.38)$$

Since $p > 2$, $C^{\frac{1}{2}-\frac{1}{p}} \subset W^{1, \frac{2p}{p+2}}$ holds. This, together with (5.33d) and (5.38), finally yields that

$$\|D_t \Psi_p(c, \gamma)\|_{L^2(r,T \times O)} \leq C_{33}(p) \left( \|\partial_t D_t \beta(c)\|_{L^2(r,T \times O)} + \|D_t c\|_{L^2(r,T \times O)} \right)$$  \hfill (5.39)

### 6 A priori estimates for SDE (2.2)

#### 6.1 Basic estimate for $y$

We begin with an $L^2$-estimate for $y$ as solution of the stochastic integral equation (4.14b). Thereby we choose $q > 2$ as in the Sobolev embedding theorem, i.e., such that $H^1(O) \subset L^q(O)$. Using assumptions (4.8) and (4.9), estimate (5.6), and a version of the Burkholder-Gundy-Davis inequality [14, Chapter 1 Theorem 7.1], we obtain that

$$\|y(t)\|_{L^q(O)} \leq C_{33} \left( \|y_0\|^q + \int_0^t a(y(s)) \|W(s)\|_{L^q(O)}^q + \int_0^t b(c,y(s)) \|s\|_{L^q(O)}^q \right)$$

$$\leq C_{33} \|y_0\|^q + t^{\frac{q-2}{2}} C_{33} \int_0^t \|a(y(s))\|_{L^q(O)}^q ds + t^{\frac{q-1}{2}} C_{33} \int_0^t \|b(c,y(s))\|_{L^q(O)}^q ds$$

$$\leq C_{33} \|y_0\|^q + C_{34} \int_0^t \|y(s)\|_{L^q(O)}^q ds + C_{34} \int_0^t \|b(y(s))\|_{L^q(O)}^q ds + \|c\|_{L^q(O)} + \|y(s)\|_{L^q(O)} \)$$

$$\leq C_{33} \|y_0\|^q + C_{35} + C_{35} \int_0^t \|y(s)\|_{L^q(O)}^q ds. \quad (6.1)$$

Integrating (6.1) over $O$ and using assumption (4.12), we conclude that

$$\|y(t)\|_{L^q(O \times (0,T) \times L^q(O))} \leq C_{33} \|y_0\|^q_{L^q(O)} + C_{36} + C_{35} \int_0^t \|y(s)\|_{L^q(O \times (0,T) \times L^q(O))}^q ds$$

$$\leq C_{37} + C_{35} \int_0^t \|y(s)\|_{L^q(O \times (0,T) \times L^q(O))}^q ds. \quad (6.2)$$

Applying the Gronwall lemma to (6.2), we arrive at the estimate

$$\|y\|_{L^q((0,T) \times L^q(O \times (0,T) \times L^q(O)))} \leq C_{38}. \quad (6.3)$$
6.2 Estimate for $\nabla y$

Computing the spatial gradient on both sides of (4.14b), we obtain that $\nabla y$ satisfies the stochastic integral equation

$$\nabla y(t) = \nabla y_0 + \int_0^t a'(y) \nabla y(s) \, dW(s) + \int_0^t \partial_c b(c, y) \nabla c(s) + \partial_y b(c, y) \nabla y(s) \, ds. \quad (6.4)$$

Using assumptions (4.8) and (4.9), estimates (5.6) and (5.7), and the Itô isometry, we obtain that

$$\|\nabla y(t)\|_{L^2(O;\Omega)^N}^2 \leq 4\|\nabla y_0\|_{L^2(O;\Omega)^N}^2 + 4\int_0^t \|a'(y)\nabla y(s)\|_{L^2(O;\Omega)^N}^2 \, ds$$

$$+ 4\int_0^t \|\partial_c b(c, y) \nabla c(s)\|_{L^2(O;\Omega)^N}^2 + \|\partial_y b(c, y) \nabla y(s)\|_{L^2(O;\Omega)^N}^2 \, ds$$

$$\leq 4\|\nabla y_0\|_{L^2(O;\Omega)^N}^2 + C_{39} \int_0^t \|\nabla c(s)\|_{L^2(O;\Omega)^N}^2 + \|\nabla y(s)\|_{L^2(O;\Omega)^N}^2 \, ds$$

$$\leq 4\|\nabla y_0\|_{L^2(O;\Omega)^N}^2 + C_{40} + C_{40} \int_0^t \|\nabla y(s)\|_{L^2(O;\Omega)^N}^2 \, ds. \quad (6.5)$$

Integrating (6.5) over $O$ and using assumption (4.12), we conclude that

$$\|\nabla y(t)\|_{L^2(O;\Omega)^N}^2 \leq 4\|\nabla y_0\|_{L^2(O;\Omega)^N}^2 + C_{41} \int_0^t \|\nabla y(s)\|_{L^2(O;\Omega)^N}^2 \, ds$$

$$\leq C_{42} + C_{40} \int_0^t \|\nabla y(s)\|_{L^2(O;\Omega)^N}^2 \, ds. \quad (6.6)$$

Applying the Gronwall lemma to (6.6), we arrive at the estimate

$$\|\nabla y\|_{L^\infty((0,T);L^2(O;\Omega))^N)} \leq C_{43}. \quad (6.7)$$

6.3 Estimate for $y(t + h_1) - y(t)$

The difference $y(t + h_1) - y(t)$ satisfies

$$y(t + h_1) - y(t) = \int_t^{t+h_1} a(y)(s) \, dW(s) + \int_t^{t+h_1} b(c, y)(s) \, ds.$$

Using assumptions (4.8) and (4.9), estimate (5.6), and a version of the Burkholder-Gundy-Davis inequality [14, Chapter 1 Theorem 7.1], we obtain that

$$\|y(t + h_1) - y(t)\|_{L^q(O)}^q \leq h_1 \frac{C_{44}}{4} \int_t^{t+h_1} \|a(y)(s)\|_{L^q(O)}^q \, ds + h_1^{q-1} C_{44} \int_t^{t+h_1} \|b(c, y)(s)\|_{L^q(O)}^q \, ds$$

$$\leq h_1 \frac{C_{45}}{4} \int_t^{t+h_1} \|y(s)\|_{L^q(O)}^q \, ds + h_1^{q-1} C_{45} \int_t^{t+h_1} \|b(s)\|_{L^q(O)}^q + \|c(s)\|_{L^q(O)}^q \, ds$$

$$\leq h_1 \frac{C_{46}}{4} + h_1^{q-1} C_{46} \int_t^{t+h_1} \|y(s)\|_{L^q(O)}^q \, ds. \quad (6.8)$$

Integrating (6.8) over $O$ and using estimate (6.3), we arrive at the estimate

$$\|y(t + h_1) - y(t)\|_{L^q(O)}^q \leq h_1^2 C_{47}. \quad (6.9)$$

6.4 Estimate for the Malliavin derivative $D_t y$

Using the chain rule and the rule of the differentiation of an Itô SDE, we compute the $D_t$-derivative on both sides of equation (4.14b). This leads to a stochastic integral equation for $D_t y$:

$$D_t y(t) = a(y)(t) + \int_r^t a'(y) D_t y(s) \, dW(s) + \int_r^t \partial_c b(c, y) D_t c(s) + \partial_y b(c, y) D_t y(s) \, ds. \quad (6.10)$$
Using assumptions (4.8) and (4.9), estimate (5.6), and the Itô isometry, we obtain that
\[
\|D_{r}y(t)\|_{L^{2}(\Omega)}^{2} \leq 4\|a(y)r\|_{L^{2}(\Omega)}^{2} + 4 \int_{r}^{t} \|a'(y)D_{r}y(s)\|_{L^{2}(\Omega)}^{2} \, ds \\
+ 4(t - r) \int_{r}^{t} \|\partial_{t} b(s, y)D_{r}c(s)\|_{L^{2}(\Omega)}^{2} + \|\partial_{y} b(s, y)D_{r}y(s)\|_{L^{2}(\Omega)}^{2} \, ds \\
\leq C_{48} \|y(r)\|_{L^{2}(\Omega)}^{2} + C_{48} \int_{r}^{t} \|D_{r}y(s)\|_{L^{2}(\Omega)}^{2} \, ds + C_{48} \int_{r}^{t} \|D_{r}c(s)\|_{L^{2}(\Omega)}^{2} \, ds.
\] (6.11)

Integrating (6.11) over \( \Omega \) and using estimate (6.3), we conclude that
\[
\|D_{r}y(t)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \leq C_{48} \|y(r)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} + C_{48} \int_{r}^{t} \|D_{r}y(s)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \, ds \\
+ C_{48} \int_{r}^{t} \|D_{r}c(s)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \, ds \\
\leq C_{48} + C_{48} \int_{r}^{t} \|D_{r}y(s)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \, ds + C_{48} \int_{r}^{t} \|D_{r}c(s)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \, ds.
\] (6.12)

Applying the Gronwall lemma to (6.12), we arrive at the estimate
\[
\|D_{r}y(t)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \leq C_{50} + C_{50} \int_{r}^{t} \|D_{r}c(s)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \, ds.
\] (6.13)

### 6.5 Estimate for \( D_{r}y(t + h_{1}) - D_{r}y(t) \)

Due to (6.10), the difference \( D_{r}y(t + h_{1}) - D_{r}y(t) \) satisfies
\[
D_{r}y(t + h_{1}) - D_{r}y(t) = \int_{t}^{t + h_{1}} a'(y(s))D_{r}y(s) \, dW(s) + \int_{t}^{t + h_{1}} \partial_{t} b(c, y)D_{r}c(s) + \partial_{y} b(c, y)D_{r}y(s) \, ds.
\]

Using assumptions (4.8) and (4.9), estimate (5.6), and the Itô isometry, we obtain that
\[
\|D_{r}y(t + h_{1}) - D_{r}y(t)\|_{L^{2}(\Omega)}^{2} \leq 3 \int_{t}^{t + h_{1}} \|a'(y(s))D_{r}y(s)\|_{L^{2}(\Omega)}^{2} \, ds + 3h_{1} \int_{t}^{t + h_{1}} \|\partial_{t} b(c, y)D_{r}c(s)\|_{L^{2}(\Omega)}^{2} + \|\partial_{y} b(c, y)D_{r}y(s)\|_{L^{2}(\Omega)}^{2} \, ds \\
\leq C_{51} \int_{t}^{t + h_{1}} \|D_{r}y(s)\|_{L^{2}(\Omega)}^{2} \, ds + h_{1} C_{51} \int_{t}^{t + h_{1}} \|D_{r}c(s)\|_{L^{2}(\Omega)}^{2} \, ds.
\] (6.14)

Integrating (6.11) over \( \Omega \) and using estimate (6.13), we conclude that
\[
\|D_{r}y(t + h_{1}) - D_{r}y(t)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \leq h_{1} C_{52} + h_{1} C_{52} \int_{t}^{t + h_{1}} \|D_{r}c(s)\|_{L^{2}(\Omega; L^{2}(\Omega))}^{2} \, ds.
\] (6.15)

### 6.6 Estimate for \( D_{r+h}y - D_{r}y \)

Due to (6.10), the difference \( D_{r+h}y - D_{r}y \) satisfies (recall that \( D_{t_{1}} F(t_{2}) = 0 \) for \( t_{2} < t_{1} \) for \( F \) adapted)
\[
D_{r+h}y(t) - D_{r}y(t) = a(y)(r + h_{2}) - a(y)(r) + \int_{r}^{t} a'(y)(D_{r+h}y - D_{r}y(s) \, dW(s) \\
+ \int_{r}^{t} \partial_{t} b(c, y)(D_{r+h}c - D_{r}c)(s) + \partial_{y} b(c, y)(D_{r+h}y - D_{r}y)(s) \, ds
\] (6.16)

Using assumptions (4.8) and (4.9), estimate (5.6), and the Itô isometry, we obtain that
\[
\|(D_{r+h}y - D_{r}y(t))\|_{L^{2}(\Omega)}^{2} \leq 4\|a(y)(r + h_{2}) - a(y)(r)\|^{2} + 4 \int_{r}^{t} \|a'(y)(D_{r+h}y - D_{r}y(s))\|_{L^{2}(\Omega)}^{2} \, ds \\
+ 4(t - r) \int_{r}^{t} \|\partial_{t} b(c, y)(D_{r+h}c - D_{r}c)(s)\|_{L^{2}(\Omega)}^{2} + \|\partial_{y} b(c, y)(D_{r+h}y - D_{r}y)(s)\|_{L^{2}(\Omega)}^{2} \, ds
\]
\[ \|y(x)\| \leq C_0, \quad \|\nabla y(x)\| \leq C_0, \quad \|\text{div} y(x)\| \leq C_0, \quad \|\Delta y(x)\| \leq C_0. \]

\[ \|\nabla^2 y(x)\| \leq C_0, \quad \|\nabla^3 y(x)\| \leq C_0, \quad \|\nabla^4 y(x)\| \leq C_0, \quad \|\nabla^5 y(x)\| \leq C_0. \]

\[ \|y(t)\| \leq C_0, \quad \|\nabla y(t)\| \leq C_0, \quad \|\text{div} y(t)\| \leq C_0, \quad \|\Delta y(t)\| \leq C_0. \]

\[ \|\nabla^2 y(t)\| \leq C_0, \quad \|\nabla^3 y(t)\| \leq C_0, \quad \|\nabla^4 y(t)\| \leq C_0, \quad \|\nabla^5 y(t)\| \leq C_0. \]

Combining (7.1) and (7.2) and using the Gronwall lemma where necessary, we arrive at the following set of estimates:

\[ \|c\| \leq C_0, \quad \|\nabla c\| \leq C_0, \quad \|\text{div} c\| \leq C_0, \quad \|\Delta c\| \leq C_0. \]

\[ \|\nabla^2 c\| \leq C_0, \quad \|\nabla^3 c\| \leq C_0, \quad \|\nabla^4 c\| \leq C_0, \quad \|\nabla^5 c\| \leq C_0. \]
and
\[|y|_{L^\infty([0,T];L^\infty(\Omega;L^4(O)))} \leq C_{17}, \]
\[|\nabla y|_{L^\infty((0,T);L^2(\Omega;L^2(O)))} \leq C_{43}, \]
\[|D_r y|_{L^\infty((r,T);L^2(\Omega;L^2(O)))} \leq C_{50}, \]
\[\left| h_2^{-\frac{1}{2}}(D_{r+h} y - D_r y) \right|_{L^\infty((r,T-h;L^2(\Omega;L^2(O))) \leq C_{56}, \]
\[\left| h_1^{-\frac{1}{2}}(D_r y(h + h_1) - D_r y) \right|_{L^\infty((r,T-h;L^2(\Omega;L^2(O)))} \leq C_{56}. \]

Estimates (7.4) allow us to apply [2, Theorem 2] directly. It yields that
\[\mathcal{Y}(T, R) \text{ is precompact in } L^2((0, T) \times O \times \Omega). \]

Remark 7.1. Observe that our estimates, particularly those involving \(D_r\), are in fact stronger than those required by that theorem. Indeed, for instance, assumptions (2)-(4) from [2, Theorem 2] deal with the regularised versions of the functions in the family.

Let us know prove the precompactness of \(\mathcal{C}(T, R)\). First, we note that (7.3f) is an estimate for a second derivative of \(\beta(c)\), not for \(c\). This precludes the direct application of [2, Theorem 2]. To overcome this, we consider instead function \(\Psi_p(c, \gamma)\). Combining the estimates (5.33) and (5.39), which we derived in Subsection 5.6, with (7.3), we obtain that
\[\left\| \nabla \Psi_p(c, \gamma) \right\|_{L^2(\Omega;L^2((0,T);L^\infty(O)))} \leq C_{27}(p), \]
\[\left\| \nabla \Psi_p(c, \gamma) \right\|_{L^2(\Omega;L^2((0,T);L^\infty(O)))} \leq C_{27}(p), \]
\[\left\| \nabla \Psi_p(c, \gamma) \right\|_{L^2(\Omega;L^2((0,T);L^\infty(O)))} \leq C_{27}(p), \]
\[\left\| D_r \Psi_p(c, \gamma) \right\|_{L^2((r,T) \times O \times \Omega)} \leq C_{57}(p), \]
\[\left\| h_2^{-\frac{1}{2}}(D_{r+h} \Psi_p(c, \gamma) - D_r \Psi_p(c, \gamma)) \right\|_{L^2((r,T-h;O \times \Omega)} \leq C_{57}(p), \]
\[\left\| D_r \Psi_p(c, \gamma) \right\|_{L^2(\Omega;C^{-\frac{1}{2}}([r,T];H^{-2}(O)))} \leq C_{57}(p). \]

With estimates (7.5) at hand we can now apply [2, Theorem 2] yielding that
\[\Psi_p(\mathcal{C}(T, R), \gamma) \text{ is precompact in } L^2((0,T) \times O \times \Omega). \]

We observe that function \(\Psi_p(\cdot, \gamma(t, x))\) has, for each fixed pair \((t, x) \in (0, T) \times O\), the following properties: it is defined on an interval, continuous, and strictly increasing. As to the latter, it is follows from the definition of \(\Psi_p\) and the fact that \(\beta'(c)\), \(\Phi_p(c, d) > 0\) for \(c, d > 0\) and \(\gamma(t, x) > 0\) for \((t, x) \in (0, T) \times O\).

Therefore, \(\Psi_p(\cdot, \gamma(t, x))\) is invertible, and its inverse has these three properties, too. Consequently, we have the following implication:

\[\Psi_p(c_n(t, x, \omega), \gamma(t, x)) \text{ is convergent in } \mathbb{R} \Rightarrow \{c_n(t, x, \omega)\}_{n \in \mathbb{N}} \text{ is convergent in } \mathbb{R}. \]

Combining (5.6) and (7.7) and using the dominated convergence theorem, we obtain that
\[\{\Psi_p(c_n, \gamma)\}_{n \in \mathbb{N}} \text{ is a.e. convergent in } (0, T) \times O \times \Omega \Rightarrow \{c_n\}_{n \in \mathbb{N}} \text{ is convergent in } L^2((0, T) \times O \times \Omega). \]

Together, (7.6) and (7.8) finally yield that
\[\mathcal{C}(T, R) \text{ is precompact in } L^2((0, T) \times O \times \Omega). \]

The proof of Theorem 4.5 is thus complete.

8 Spatial semi-discretization for a nondegenerate case

In this section we set up and study a spatial finite-difference scheme for system (2.1)-(2.2) under the additional assumption
\[\beta \in C^2([0, \mathbb{R}^+]) \] which corresponds to a nondegenerate case. Our goal here is twofold. On one hand, we apply the semi-discretisation method (see, e.g. [12, Chapter 4 §1] and references therein, particularly [18]) in order to obtain the existence of solutions. At the same time, we illustrate thereby how one can use compactness in order to rigorously prove the convergence of a numerical scheme for nonlinear systems with stochasticity.
8.1 Spatial discretisation and interpolation

In this sequel we recall some concepts and ideas of the deterministic semi-discretisation method. We refrain from the proofs of the properties listed here since they either exactly repeat or are slight modifications of results addressed in the literature. The interested reader is referred to [12, Chapter 4 §1], where the method is described. The discretization is preformed only in $O$. To avoid some purely deterministic technical difficulties which have to do with discretizing close to and on the boundary of $O$, we restrict our exposition to the case when the spatial domain $O$ is a unit cube:

$$O = (0,1)^N.$$  

However, we especially emphasise that this simplification is by no means essential: the present approach can be used for more general domains.

We begin with some more notation. As usual, we denote by $e_k$ the $k$-th standard basis vector in $\mathbb{R}^N$. Let $M \in \mathbb{N}\backslash\{1\}$. For $h := \frac{1}{M+1}$ we define the discrete sets

$$O^h := \{h,2h,\ldots,1-h\}^N, \quad \overline{O}^h := \{0,h,\ldots,1\}^N, \quad \partial^h O^h := \overline{O}^h \backslash O^h$$

and the space-continuous set

$$U^h := (2h,1-2h)^N.$$  

Clearly, it holds that

$$|O \backslash U^h| = 1 - (1 - 4h)^N \to 0.$$  

Next, we introduce standard finite difference operators for a function $u$ at a point $m$:

$$\partial^h_k u(m) := \frac{1}{h}(u(m+he_k) - u(m)) \quad \text{for } k \in \{1,\ldots,N\}, \quad (8.2)$$

$$\nabla^h := (\partial^h_1, \ldots, \partial^h_N)^T, \quad (8.3)$$

$$\Delta^h u(m) := \frac{1}{h^N} \sum_{k=1}^N (u(m+he_k) - 2u(m) + u(m-he_k)), \quad (8.4)$$

$$\partial^h_{\infty} u(m) := \frac{1}{h}(u(m) - u(\overline{m})) \quad \text{for } m \in \partial^h O^h \quad (8.5)$$

where

$$\overline{m}_k := \begin{cases} h & \text{if } m_k = 0, \\ m_k & \text{if } m_k \in (0,1), \\ 1-h & \text{if } m_k = 1. \end{cases} \quad (8.6)$$

We recall the discrete version of the Green’s first identity (i.e., summation by parts formula)

$$\sum_{i=1}^M (a_{i+1} - 2a_i + a_{i-1})b_i = -\sum_{i=0}^M (a_{i+1} - a_i)(b_{i+1} - b_i) + (a_{M+1} - a_M)b_{M+1} - (a_1 - a_0)b_0, \quad (8.7)$$

as well as the following relation

$$\partial^h_k (g(u))(m) = \partial^h_k u(m) \int_0^1 g'(u(m) + \tau h \partial^h_k u(m)) \, d\tau$$ \quad (8.8)

which serves as a sort of chain rule for the discrete case. Further, we make use of some discrete analogs of several (semi-)norms and a scalar product which mimic the corresponding notions in connection with the Lebesgue and Sobolev spaces: for $u,v : \overline{O}^h \to \mathbb{R}$ let

$$(u,v)_{L^2(D^h,h)} := h^N \sum_{m \in D^h} uv(m) \quad \text{for } D^h \subset \overline{O}^h,$$

$$\|u\|_{L^p(D^h,h)} := \begin{cases} (h^N \sum_{m \in D^h} |u(m)|^p)^{\frac{1}{p}} & \text{for } p \in [1,\infty), \\ \max_{m \in D^h} |u(m)| & \text{for } p = \infty, \end{cases} \quad \text{for } D^h \subset \overline{O}^h,$$
Moreover, both interpolation operators preserve positivity:

\[ |u|_{H^1(\Omega; \mathcal{h})} := \left( \frac{1}{h^{N-2}} \sum_{m, m+h \in \mathcal{G}} |u(m + h e_k) - u(m)|^2 \right)^{\frac{1}{2}}, \]

\[ \|u\|_{H^1(\Omega; \mathcal{h})} := \left( |u|_{H^1(\Omega; \mathcal{h})}^2 + |u|_{L^2(\Omega; \mathcal{h})}^2 \right)^{\frac{1}{2}}, \]

\[ \|u\|_{H^{2-}(\Omega; \mathcal{h})} := \sup_{v \in H^2(\Omega; \mathcal{h}) \setminus \{0\}} \frac{|(u, v)_{L^2(\Omega; \mathcal{h})}|}{\|\Delta v\|_{L^2(\Omega; \mathcal{h})}}. \]

where

\[ H^2_0(\Omega; \mathcal{h}) := \{ v : \Omega^h \rightarrow \mathbb{R} : v = \partial^h v = 0 \text{ on } \partial^h \Omega \}. \]

For a function \( u : \Omega^h \rightarrow \mathbb{R} \) we define its projection to the space of discrete functions via

\[ P^h u(m) := \int_{(m + (-\frac{1}{2}, \frac{1}{2})) \cap \Omega} u(z) \, dz \quad \text{for } m \in \Omega^h. \]

In order to interpolate discrete functions \( u : \Omega^h \rightarrow \mathbb{R} \) we use two types of splines: the piecewise constant

\[ \Pi^h u(x) := \begin{cases} u(m) & \text{for } x \in (m + (-\frac{1}{2}, \frac{1}{2})) \cap \Omega, \ m \in \Omega_h, \\ 0 & \text{otherwise} \end{cases} \]

and the piecewise polyaffine

\[ \Lambda^h u(x) := \sum_{i \in (m + (0, h)) \cap \Omega_h} u(i) L^h_+(i-m) \left( \frac{1}{h}(x-m) \right) \quad \text{for } x \in (m + [0, h]^N) \cap \Omega, \ m \in \Omega_h, \]

where

\[ \Delta^h(z) := \prod_{j=1}^N \beta_{l_j}(z_j) \quad \text{for } z \in [0, 1]^N, \ l \in \{0, 1\}^N, \]

\[ \beta_0(r) := 1 - r, \ \beta_1(r) := r \quad \text{for } r \in [0, 1]. \]

The constructed splines have the following useful properties: for all \( u : \Omega^h \rightarrow \mathbb{R} \) it holds that

\[ \Pi^h u \in L^\infty(\Omega), \]

\[ (\Pi^h u, \Pi^h v)_{L^2(\Omega; \mathcal{h})} = (u, v)_{L^2(\Omega; \mathcal{h})} \quad \text{for } v \in H^2_0(\Omega; \mathcal{h}), \] (8.9)

\[ \varphi(\Pi^h u) = \Pi^h \varphi(u) \quad \text{a.e. in } \Omega \text{ for any } \varphi : \mathbb{R} \rightarrow \mathbb{R}, \] (8.10)

\[ \Lambda^h u \in W^{1, \infty}(\Omega), \]

\[ \Lambda^h u(m) = u(m) \quad \text{for } m \in \Omega^h, \]

and

\[ u = 0 \quad \text{in } \partial^h \Omega, \quad \Rightarrow \quad \Lambda^h u = 0 \quad \text{in } \partial \Omega, \] (8.11)

\[ \partial^h u = 0 \quad \text{in } \partial^h \Omega, \quad \Rightarrow \quad \partial^h \Lambda^h u = 0 \quad \text{a.e. in } \partial \Omega. \] (8.12)

Moreover, both interpolation operators preserve positivity:

\[ u \geq 0 \quad \text{in } \Omega^h, \quad \Rightarrow \quad \Lambda^h u, \Pi^h u \geq 0 \quad \text{in } \Omega. \] (8.13)

Finally, the projection and interpolations enjoy the following estimates:

\[ \left| P^h u \right|_{L^p(\Omega; \mathcal{h})} \leq C \| u \|_{L^p(\Omega)} \quad \text{for } u \in L^p(\Omega), \ p \in [1, \infty], \] (8.14)

\[ \left| P^h u \right|_{H^1(\Omega; \mathcal{h})} \leq C \| u \|_{L^2(\Omega)} \quad \text{for } u \in H^1(\Omega), \] (8.15)
where the spatially discretized initial data 

\[ \text{systems with smooth coefficients} \]

The nonnegativity of solution components is a consequence of a general result on the invariance for SDE.

\[ R(\beta) \]

functions

A solution can then be understood in the usual strong SDE-sense. Observe that not only the coefficient \( v \)

In particular, one can take \( v^h := v^h_{\Omega^h} \) for any \( v \in C_0^\infty(U^h) \) as a test function in (8.25). Thus, using the interpolation operators and properties (8.9), (8.10), and (8.11) ((8.12)) one deduces from the space-

8.2 Approximation via semi-discretization

We start with the following semi-discretization of the RPDE-SDE system (2.1)-(2.2):

\[
\begin{aligned}
\frac{d}{dt} c^h &= \beta'(c^h) \Delta^h c^h + f(c^h, y^h) \\
\frac{d}{dt} y^h &= a(y^h) \, dW^h + b(c^h, y^h) \, dt
\end{aligned}
\]

in \( \Omega^h \) \((8.22a)\)

\[
\begin{aligned}
\frac{d}{dt} c^h &= 0 \\
\frac{d}{dt} y^h &= y_0^h
\end{aligned}
\]

in \( \Omega \) \((8.22b)\)

where the spatially discretized initial data \((c_0^h, y_0^h) : \Omega^h \to \mathbb{R}_0^+\) is defined via

\[
(c_0^h, y_0^h) := (P^h c_0, P^h y_0) \quad \text{in} \quad \Omega^h.
\]

Observe that equation (8.22a) can be rewritten in the conventional form of a RODE:

\[
\frac{d}{dt} c^h = (\beta'(c^h))^{-1} \Delta^h c^h + (\beta'(c^h))^{-1} f(c^h, y^h).
\]

Using the boundary conditions (8.22b) one can eliminate \((c^h(\cdot, m, \cdot))_{m \in \mathbb{N}}\) from (8.22a). Then, system (8.22)-(8.23) can be considered as a RODE-SDE system with respect to

\[
((c^h(\cdot, m, \cdot))_{m \in \mathbb{N}}, (y^h(\cdot, m, \cdot))_{m \in \mathbb{N}}) : [0, T] \times \Omega \to (\mathbb{R}_0^+)^{M+1} \times (\mathbb{R}_0^+)^{M+1}.
\]

A solution can then be understood in the usual strong SDE-sense. Observe that not only the coefficient functions \( f, a, b \), but also \((\beta')^{-1} \) and \( \Delta^h : \Omega^h \to \mathbb{R}_0^+ \) in equations (8.24) and (8.23a) are continuously differentiable in \( \mathbb{R}_0^+ \). For \( \beta \) this holds since \((\beta')^{-1}, \beta'' \in C(\mathbb{R}_0^+), \) the latter due to our additional assumption (8.1), while \( \Delta^h \) is simply linear. Therefore, the RODE-SDE is uniquely solvable and its solutions possess square integrable Malliavin derivatives due to well-known results [16, Corollary 2.2.1, Theorem 2.2.1].

The nonnegativity of solution components is a consequence of a general result on the invariance for SDE systems with smooth coefficients [15] and assumptions (4.10) and (4.13).

Multiplying (8.22a) by \( \xi^h \) for arbitrary \( \xi \in C_0^\infty(0, T) \) and \( v^h \in H_0^2(\Omega^h : h) \), summing over \( \Omega^h \) thereby using (8.7) twice, and integrating over \((0, T)\) we obtain the following weak formulation:

\[
\begin{aligned}
- (\beta(c_0^h), v^h)_{L^2(\Omega^h : h)} \xi(0) - \int_0^T (\beta(c^h)(s), v^h)_{L^2(\Omega^h : h)} \frac{d\xi}{dt}(s) \, ds &= \\
\int_0^T (\Delta^h v^h)(s, \Omega^h) \xi(s) + (f(c^h, y^h)(s), v^h)_{L^2(\Omega^h : h)} \xi(s) \, ds.
\end{aligned}
\]

In particular, one can take \( v^h := v^h_{\Omega^h} \) for any \( v \in C_0^\infty(U^h) \) as a test function in (8.25). Thus, using the interpolation operators and properties (8.9), (8.10), and (8.11) ((8.12)) one deduces from the space-
discrete problem (8.22)-(8.23) an approximation to the original space-continuous system (2.1)-(2.2):

\[
\begin{align*}
- (\beta (\Pi^h P^h c_0), \Pi^h v)_{L^2(O)} & \xi (0) - \int_0^T (\beta (\Pi^h c^h) (s), \Pi^h v)_{L^2(O)} \frac{d^x}{dt} (s) \, ds \\
= & \int_0^T (\Pi^h c^h(s), \Pi^h \Delta^h v)_{L^2(O)} \xi(s) + (f (\Pi^h c^h, \Pi^h y^h) (s), \Pi^h v)_{L^2(O)} \xi(s) \, ds,
\end{align*}
\]

for all \( v \in C^\infty_c(U^h) \), \( \xi \in C^\infty_0 [0, T) \) a.s. in \( \Omega \),

\[
\Lambda^h c^h = 0 \quad \text{in } (0, T) \times \partial O \\
(\partial_t \Lambda^h c^h = 0 \quad \text{a.e. in } (0, T) \times \partial O) \quad \text{a.s. in } \Omega, \tag{8.26a}
\]

\[
\Pi^h y^h(s) = \Pi^h P^h y_0(t) + \int_0^t \Pi^h y^h(s) \, dW(s) + \int_0^t b (\Pi^h c^h, \Pi^h y^h) (s) \, ds
\]
in \( L^2(O) \) for all \( t \in [0, T] \) a.s. in \( \Omega \). \tag{8.26c}

This system will be analysed in the subsequent subsections.

### 8.3 A priori estimates

It is easy to see that for solutions of the semi-discrete system (8.22)-(8.23) one has a set of uniform in \( h \) estimates which is very similar to estimates (7.3)-(7.4) for the original space-continuous system (2.1)-(2.2). Indeed, one just needs to replace the spatial differential operators by their discrete versions (8.2)-(8.6), integration over \( O \) by summation over \( O^h \), and the \( L^p(O) \) norms by their discrete analogues \( L^p(O^h; h) \) and to use (8.14)-(8.15) in order to estimate \( (c^h_0, y^h_0) \), the discrete version of the Green’s first identity (8.7), the boundary conditions (8.22b), as well as relation (8.8) instead of the chain rule while dealing with spatial derivatives. Moreover, since we assumed that \( \beta \in C^1 (\mathbb{R}^+ \times \mathbb{R}^+) \), the uniform estimates for \( c^h \) can be directly transformed into the corresponding uniform estimates for \( \beta (c^h) \). This spares the need of constructing more complicated transformations such as those derived in Subsections 5.5.5-5.5.6. We thus get the following sets of estimates:

\[
\| \beta (c^h) \|_{L^\infty(\Omega, L^\infty((0,T);L^\infty(\overline{\Omega};h)))} \leq C_{59}, \tag{8.27a}
\]

\[
\| \beta (c^h) \|_{L^\infty(\Omega, L^\infty((0,T);H^1(\overline{\Omega};h)))} \leq C_{59}, \tag{8.27b}
\]

\[
\| \partial_t \beta (c^h) \|_{L^\infty(\Omega, L^2((0,T);L^2(\overline{\Omega};h)))} \leq C_{59}, \tag{8.27c}
\]

\[
\| D_r \beta (c^h) \|_{L^2((r,T,\Omega;L^2(\overline{\Omega};h)))} \leq C_{59}, \tag{8.27d}
\]

\[
\| h_2^{\frac{3}{2}} (D_{r+h_2} \beta (c^h) - D_r \beta (c^h)) \|_{L^2((r,T-h_2,\Omega;L^2(\overline{\Omega};h)))} \leq C_{59}, \tag{8.27e}
\]

\[
\| \partial_t D_r \beta (c^h) \|_{L^2((r,T,\Omega;L^2(\overline{\Omega};h)))} \leq C_{59} \tag{8.27f}
\]

and

\[
\| y^h \|_{C^\frac{1}{2}((0,T);L^\infty(\Omega, L^\infty(\overline{\Omega};h)))} \leq C_{59}, \tag{8.28a}
\]

\[
\| y^h \|_{L^\infty((0,T);L^2(\Omega;H^1(\overline{\Omega};h)))} \leq C_{59}, \tag{8.28b}
\]

\[
\| D_r y^h \|_{L^\infty((r,T,\Omega;L^2(\overline{\Omega};h)))} \leq C_{59}, \tag{8.28c}
\]

\[
\| h_2^{\frac{3}{2}} (D_{r+h_2} y^h - D_r y^h) \|_{L^\infty((r,T-h_2,\Omega;L^2(\overline{\Omega};h)))} \leq C_{59} \tag{8.28d}
\]

\[
\| h_1^{\frac{1}{2}} (D_r y^h (\cdot + h_1) - D_r y^h) \|_{L^\infty((r,T-h_1,\Omega;L^2(\overline{\Omega};h)))} \leq C_{59} \tag{8.28e}
\]

Combining (8.27)-(8.28) with (8.17)-(8.19) we conclude that the pair \( (\Lambda^h \beta (c^h), \Lambda^h (y^h)) \) satisfies

\[
\| \Lambda^h \beta (c^h) \|_{L^\infty((0,T);L^\infty(\Omega, L^\infty(O)))} \leq C_{60}, \tag{8.29a}
\]

\[
\| \Lambda^h \beta (c^h) \|_{L^\infty((0,T);H^1(\Omega, h)))} \leq C_{60}, \tag{8.29b}
\]

\[
\| \partial_t \Lambda^h \beta (c^h) \|_{L^\infty((0,T);L^2(\Omega, L^2(O)))} \leq C_{60}, \tag{8.29c}
\]

\[
\| D_r \Lambda^h \beta (c^h) \|_{L^2((r,T;L^2(\Omega, L^2(O))))} \leq C_{60}, \tag{8.29d}
\]

\[
\| h_2^{\frac{3}{2}} (D_{r+h_2} \Lambda^h \beta (c^h) - D_r \Lambda^h \beta (c^h)) \|_{L^2((r,T-h_2,\Omega;L^2(O)))} \leq C_{60}, \tag{8.29e}
\]

\[
\| h_1^{\frac{1}{2}} (D_r \Lambda^h \beta (c^h) (\cdot + h_1) - D_r \Lambda^h \beta (c^h)) \|_{L^2((r,T-h_1,\Omega;L^2(O)))} \leq C_{60}. \tag{8.29f}
\]
\[ \left\| \partial_t D_r \Lambda^h \beta (c^h) \right\|_{L^2((r,T) \times \Omega; L^2(U^h_n)))} \leq C_{60}, \]  
\text{(8.29f)}

\[ \left\| \Lambda^h y^h \right\|_{L^2([0,T]; L^2(\Omega; L^2(U^h_n)))} \leq C_{60}, \]  
\text{(8.30a)}
\[ \left\| \Lambda^h y^h \right\|_{L^\infty((0,T); L^2(\Omega; H^1(\Omega)))} \leq C_{60}, \]  
\text{(8.30b)}
\[ \left\| D_r \Lambda^h y^h \right\|_{L^\infty((0,T); L^2(\Omega; L^2(U^h_n)))} \leq C_{60}, \]  
\text{(8.30c)}
\[ \left\| h^2 \left( D_{r+h_2} \Lambda^h y^h - D_r \Lambda^h y^h \right) \right\|_{L^\infty((r,T-h_2: L^2(\Omega; L^2(U^h_n)))} \leq C_{60} \]  
\text{(8.30d)}
\[ \left\| h^1 \left( D_r \Lambda^h y^h (\cdot + h_1) - D_r \Lambda^h y^h (\cdot) \right) \right\|_{L^\infty((r,T-h_2: L^2(\Omega; L^2(U^h_n)))} \leq C_{60}. \]  
\text{(8.30e)}

Further, (8.27b) and (8.28b) together with (8.20) and assumption \( \beta^{-1} \in C^1(\mathbb{R}_0^+) \) (compare (4.4)), yield that

\[ \left\| \Pi^h y^h - \Lambda^h y^h \right\|_{L^\infty((0,T); L^2(\Omega; L^2(U^h_n)))} \leq hC_{61}. \]  
\text{(8.31)}
\[ \left\| \Pi^h \beta (c^h) - \Lambda^h \beta (c^h) \right\|_{L^\infty((0,T); L^2(\Omega; L^2(U^h_n)))} \leq hC_{61} \]  
\text{(8.32)}

and

\[ \left\| \Pi^h y^h - \Lambda^h y^h \right\|_{L^\infty((0,T); L^2(\Omega; L^2(U^h_n)))} \leq hC_{61}. \]  
\text{(8.33)}

Finally, due to (8.16) we also have for the approximations of the initial data the estimates

\[ \left\| \Pi^h y^h_0 - c_0 \right\|_{L^2(\Omega)} \leq hC_{62}, \]  
\text{(8.34)}
\[ \left\| \Pi^h y^h_0 - y_0 \right\|_{L^2(\Omega)} \leq hC_{62}. \]  
\text{(8.35)}

### 8.4 Compactness and convergence

Thanks to estimates (8.29)-(8.30) we are once again in the position when we can directly apply [2, Theorem 2]. It yields that

\[ \left\{ \left( \Lambda^h (c^h), \Lambda^h y^h \right) : \ h = \frac{1}{M+1}, \ M \in \mathbb{N} \right\} \text{ is precompact in } \left\{ L^2((0, T) \times \Omega ; \mathbb{R}^2) \right\}. \]  
\text{(8.36)}

Together with estimates (8.32)-(8.33) this leads to

\[ \left\{ \left( \Pi^h \beta (c^h), \Pi^h y^h \right) : \ h = \frac{1}{M+1}, \ M \in \mathbb{N} \right\} \text{ is precompact in } \left\{ L^2((0, T) \times \Omega ; \mathbb{R}^2) \right\}. \]  
\text{(8.37)}

Further, using \( \beta^{-1} \in C^1(\mathbb{R}_0^+) \) and property (8.10) for \( \varphi := \beta \), we conclude with (8.37) that

\[ \left\{ \left( \Pi^h c^h, \Pi^h y^h \right) : \ h = \frac{1}{M+1}, \ M \in \mathbb{N} \right\} \text{ is precompact in } \left\{ L^2((0, T) \times \Omega ; \mathbb{R}^2) \right\}. \]  
\text{(8.38)}

Combining estimates (8.29)-(8.30), (8.31), (8.33) with the compactness results (8.36) and (8.38) and using assumption \( \beta^{-1} \in C^1(\mathbb{R}_0^+) \) we deduce that there exists a pair of functions \( (c, y) : [0,T] \times \overline{\Omega} \times \Omega \rightarrow \mathbb{R}_0^+ \times \mathbb{R}_0^+ \) which satisfies conditions 1-6 from Definition 4.2 and a sequence \( h_n \rightarrow 0 \) such that

\[ \left( \Pi^h c^h_n, \Pi^h y^h_n \right) \rightharpoonup (c, y) \text{ in } L^2((0,T) \times \Omega), \text{ in } L^2((0,T) \times \Omega) \text{ a.s. in } \Omega \]  
\text{(8.39)}

and

\[ \Lambda^h_n c^h_n \rightharpoonup c \text{ in } L^\infty(\Omega; L^\infty((0,T); H^1(\Omega))). \]  
\text{(8.40)}

Due to (8.34)-(8.35) it holds also that

\[ \left( \Pi^h y^h_0, \Pi^h y^h_0 \right) \rightharpoonup (c_0, y_0) \text{ in } (L^2(\Omega))^2. \]  
\text{(8.41)}
Using (8.39) and (8.41), assumption $\beta, f, a, b \in C^1$, and the fact that
\[ \Pi^h \Delta^h v \xrightarrow{n \to \infty} \Delta v \quad \text{in } L^2(O) \]
due to (8.21), we can pass to the limit in the weak formulation (8.26a) and the SDE (8.26c). Finally, combining (8.40) with the boundary conditions (8.26b) and the trace theorem, we conclude that $c$ satisfies either of the boundary conditions. Altogether, this means that $(c, v)$ is a weak solution to the PDE (2.1) and a strong solution to the SDE (2.2), and it has the desired regularity. Thus, we proved the following Lemma:

**Lemma 8.1** (Existence for a nondegenerate case). Let Assumptions 4.1 be satisfied. Assume in addition that $\beta \in C^2(\mathbb{R}_0^+)$ Then there exists a weak-strong global solution in terms of Definition 4.2 to system (2.1)-(2.2).

9 Proof of the existence Theorem 4.6

In this final section we discuss how our compactness Theorem 4.5 can be used in order to prove the existence of solutions to the original degenerate system (2.1)-(2.2). Let $T > 0$ be arbitrary. Choose $R_0, R_1 > 0$ large enough so that $(\beta, f, a, b, c_0, y_0) \in \mathcal{P}(T, (R_0/2, R_1/2))$. For each $\varepsilon \in (0, 1)$ let $\beta_\varepsilon$ be such that
\[
\beta_\varepsilon, f, a, b, c_0, y_0) \in \mathcal{P}(T, (R_0, R_1)),
\]
\[
\beta_\varepsilon \in C^2(\mathbb{R}_0^+),
\]
\[
\beta_\varepsilon \xrightarrow{\varepsilon \to 0} \beta \quad \text{in } C[0, R_2],
\]
where $R_2$ is an upper bound for the $L^\infty$-norm of $c$, compare (4.15). We introduce a family of approximating problems
\[
\begin{cases}
\partial_t \beta_\varepsilon(c_\varepsilon) = \Delta c_\varepsilon + f(c_\varepsilon, y_\varepsilon) & \text{in } (0, T] \times O, \\
\beta_\varepsilon = 0 & \text{in } (0, T] \times \partial O, \\
c_\varepsilon = c_0 & \text{in } \{0\} \times O.
\end{cases}
\]
\[
\begin{cases}
dy_\varepsilon = a(y_\varepsilon) \, dW + b(c_\varepsilon, y_\varepsilon) \, dt & \text{in } (0, T] \times O, \\
y_\varepsilon = y_0 & \text{in } \{0\} \times O.
\end{cases}
\]

The existence of a solution $(c_\varepsilon, y_\varepsilon)$ in terms of Definition 4.2 to system (9.2)-(9.3) follows from Lemma 8.1 of the previous section. Moreover, due to assumption (9.1) and the compactness Theorem 4.5 family \{$(c_\varepsilon, y_\varepsilon)$\}, $\varepsilon \in (0, 1)$ is precompact in $(L^2((0, T) \times O \times O))^2$. Standard compactness arguments and a limit procedure then yield the existence of a solution in terms of Definition 4.2 to the original degenerate system (2.1)-(2.2) for any $T > 0$.

Finally, we observe that the particular choice of the zero starting time was not essential for our arguments heretofore. Indeed, all previous results continue to hold if we replace the interval $[0, T]$ by $[t_0, t_0 + T]$ for any $t_0 > 0$ and consider the ‘shifted’ in time filtration $(\mathcal{F}_{t+t_0})_{t \geq 0}$ instead of the original one. Thus, we obtain a solution defined for all times by defining it successively in $[0, T], [T, 2T]$, and so on. The proof of Theorem 4.6 is thus complete.

**Remark 9.1** (Numerical schema). As a by-product of the constructions from this Section and the previous Section 8 we have a numerical schema for (2.1)-(2.2) which is, at least theoretically, converging.

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