A Novel Method to Design Stub-Loaded Microstrip Filters with Arbitrary Passband Based on the Reflection Theory and Monte Carlo Method

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Abstract. In this paper, a novel method to design microstrip filters with arbitrary passband is proposed. A kind of stub-loaded microstrip filter with asymmetric structure is proposed in this paper. The mathematic function of the total reflection coefficient at the input port of the filter is derived based on the reflection theory, and an ultra-wide-band (UWB) filter which can work from 11 GHz to 39 GHz is fabricated and measured to verify that our function can calculate the reflection coefficient precisely. The circuit model of this UWB filter is provided, and the comparison between frequency responses of the circuit simulation and the calculation results justifies the validity of our proposed method. The procedure based on the Monte Carlo method to design stub-loaded microstrip filters with arbitrary passbands is given in this paper. A bandpass filter (BPF) which can work on 7–9.5 GHz is designed and measured as an example. The comparison between calculated dimensions and practical dimensions of this filter shows a satisfactory fitting degree. Then the proposed method is compared with the Chebyshev filter design method to show the advantages of our proposed method. A low-pass filter (LPF) which can work below 4 GHz is designed and compared with some prior arts in order to show that our proposed filter is simple in structure and small in size. At last, the adaptabilities of our proposed method in narrow band and multi-band situations are discussed in this paper. By using our proposed method, the stub-loaded microstrip filter can be designed quickly and easily. Based on the results mentioned above, we can count that the method proposed in this paper has a good application in microstrip filter design.

Keywords
Electromagnetic reflection, microstrip filters, Monte Carlo methods, reflection coefficient

1. Introduction

1.1 Motivation and Related Works

Microstrip filters have been widely used in a large number of RF applications. As the wireless communication systems are developing rapidly today, it is important for RF engineers to hunt for some convenient methods to design microstrip filters [1].

A general design method for multiband filters is proposed in [2]. In [2], the frequency and element transformations from a low-pass prototype filter to a practical multiband bandpass filter is derived, and the formula which extracts the coupling coefficient \( k \) between coupled multimode resonators and the external quality factor are obtained. A novel method to design low loss dual-band microstrip filters using folded open-loop ring resonators (OLRRs) is presented in [3]. The first and second passbands of the dual-band filter can be adjusted to a desired frequency bands by adjusting the physical dimensions of OLRRs. A novel microstrip Chebyshev low-pass filter (LPF) design method using a defected ground structure (DGS) is introduced in [4]. The DGS reduces the lengths of inductive microstrip lines effectively, so the miniaturization of the filter dimensions is achieved. In this paper, a novel method to design stub-loaded microstrip filters is proposed. The whole design procedure is based on the reflection theory and Monte Carlo method.

In [5], the theory of small reflections (TSR) and its derivation process is introduced. The TSR is to consider the total reflection coefficient at the input port as the composite function of local reflection coefficients between the neighboring transmission lines [5], [6]. This theory is widely used in the design of impedance matching between unmatched sources and loads [7–11]. In [9], a microstrip feeding for an antenna is designed based on the TSR, and the return loss in its working frequency band is under \(-60\ dB\). This result verifies that the TSR has good guidance in impedance matching design.
Stepped-impedance filters also apply the TSR to design some well performed low-pass filters (LPFs) [12–17]. In [17], the authors generalize this theory for the case of multi-section transformers. In [24], a LPF which can work below 4 GHz is designed based on the stepped-impedance microstrip lines, and it is an example for the traditional design procedure of stepped-impedance microstrip LPFs.

The filters applying stepped-impedance resonators are studied in recent years [18–23]. A microstrip low-pass filter (LPF) with three coupled stepped impedance resonators is presented in [21]. This filter has wide stopband bandwidth and can offer a high selectivity. In [22], a wide-stopband bandpass filter based on slitted stepped-impedance resonators is introduced. Using slitted stepped-impedance resonators can widen the bandwidth of the stop-band of the filter and add more freedoms in controlling resonance modes compared with a stepped-impedance resonator. Stepped-impedance ring resonator is also used to construct balanced tri-band bandpass filters which have excellent out-of-band performance [23].

However, there is a limitation while using the TSR to design microstrip filters: if the impedance between two neighboring stubs has a large discrepancy, some approximations based on the TSR will be incorrect. In [25], the authors present a modified version of TSR and use this technology to design a broadband filter, but the frequency response of this filter dose not perform well.

In [26], we propose a method based on the TSR to design microstrip filters with arbitrary passbands. In that paper, a novel stub-loaded microstrip filter is proposed and the approximative reflection coefficient function at the input port of the filter is derived. In this paper, we present a modified version of the method in [26]. Our modified method can achieve more accurate calculation results than the method in [26]. The reflection coefficient function becomes very complex, so it is difficult for us to solve the dimensions of our filters from the reflection coefficient function. Therefore, it is important for us to find some feasible calculation methods to figure out this problem.

In this paper, we choose the Monte Carlo method as our reverse calculation method. The Monte Carlo method is an efficient mathematical tool to solve some extremely complicated problems [27]. In [28], the authors propose a transition between a microstrip line and a ridge gap waveguide, and they use the Monte Carlo method to analysis the assembly tolerances.

By using our modified method, the design procedure is simplified and the proposed microstrip filter can be designed simply and quickly compared with other methods.

### 1.2 Summary of the Contributions and Paper Organization

The main contributions of our present study are summarized as follows:

- A novel multiple-stub-loaded impedance transformer filter is proposed in this paper. We use the reflection theory to calculate the total reflection coefficient of this types of filter and derive the accurate reflection function at the input port. To verify that our proposed function works well, an UWB filter is simulated and fabricated, and the calculation results of S-parameter fit well with the practical measurement results. Also, the circuit model of this UWB filter is given, and the frequency responses of the circuit simulation are compared with the calculation results of the proposed method. The comparison results verify that our proposed method can work as good as the circuit method.

- The proposed reflection function is used to calculate the practical dimensions of our filter. To achieve this purpose, the Monte Carlo method is used as the reverse solution method to find the optimal approximation of our design goals. The relationship between the number of Monte Carlo experiments and the fitting error is investigated. The whole procedure is given in the paper and a BPF is designed by our method to verify the feasibility of our proposed method. In Sec. 3, the proposed BPF is redesigned by the Chebyshev method. The design procedure of the Chebyshev method is compared with our proposed calculation procedure, and the results show that our method can obtain the dimensions of the filter simply and directly.

- In Sec. 3, the comparisons between our proposed method and other prior arts are shown, and the characteristics of our proposed method are discussed. A LPF is designed and compared with filters proposed in [24] and [29] in order to show the advantages of our proposed method and filter. At last, the frequency responses of a dual-band microstrip filter proposed in [26] is recalcualted by the method proposed in this paper to justify the application of our modified method in the narrow band and multi-band situations.

The rest of this paper is organized as follows. Section 2 shows our theory and contains the derivation of the total reflection coefficient of multiple-stub-loaded filters at the input port. Section 3 shows the design procedure which bases on the Monte Carlo experiments, and the advantages of our proposed method are also discussed in Sec. 3. Finally, we conclude our paper in Sec. 4.

There are lots of notations used in this paper and they are summarized in Tab. 1.
2. Theory Background

2.1 Problem Statement

The [5] shows the method which is used to obtain the total reflection coefficient $|\Gamma_{\text{total}}|$ of the multiple-stub impedance transformer at the input port. This method is known as the TSR. In [26], a modified version of this theory is proposed and implemented in microstrip filter design, and a novel kind of sub-loaded filters is proposed. The parallel opened- and shorted-stubs are used to change the norm of the local reflection coefficients of the main transmission line in this novel stub-loaded filter. (The detail of this operation is explained in (5)). By using this method, the reflection curve can be shaped in some more complex forms [26].

The equivalent model of this type of filter is shown in Fig. 1. Where $L$ represents the length of each part of the main transmission line, $L_{x}^{(n)}$ represents the dimension of each parallel loaded-stub, $N$ represents the orders of the filter, and $\Gamma_{n}$ represents the local reflection coefficient of neighboring parts.

The $|\Gamma_{\text{total}}|$ function is shown in (3), where $f$ represents the frequency:

$$Y_{n} = Y_{o} + Y_{\text{im}}^{(n)}$$

$$\Gamma_{n}(f) = \frac{Y_{o}^{(n)}(L_{x}^{(n)}, f) - Y_{n+1}^{(n)}(L_{x}^{(n+1)}, f)}{Y_{o}^{(n)}(L_{x}^{(n)}, f) + Y_{n+1}^{(n)}(L_{x}^{(n+1)}, f)}.$$  \hspace{1cm} (2)

$$|\Gamma_{\text{total}}(f)| = \sum_{N} |\Gamma_{n}(f)| \exp(-\frac{4\pi nL}{c} f).$$  \hspace{1cm} (3)

In [26], equation (3) is used to design the frequency response of the filter. Two microstrip filters are designed and measured. The measured results fit well with the calculated results. But there are some limits while using the method proposed in [26]:

- Equation (3) is based on the TSR, and it is an approximation formula of the practical $|\Gamma_{\text{total}}|$ function. When the characteristic impedances between the neighbouring parts of the filter differ sharply, this theory will not be applicable. Therefore, the calculation results of the $|\Gamma_{\text{total}}|$ curve can only indicate the tendency of the practical $|\Gamma_{\text{total}}|$ curve but cannot calculate the curve precisely.

- Equation (3) is simplified and shown in (4) while the structure of the filter is chosen to use the centro-symmetric structure.

\[
|\Gamma_{\text{total}}(f)| = 2 |f| \exp(-\frac{24\pi nL}{c} f) \times \sum_{N} |\Gamma_{n}(L_{n}, L_{n+1}, f)| \sin\left(\frac{4\pi nL}{c} f\right). \hspace{1cm} (4)
\]

This means that a pair of symmetric stubs can only make contributions to one $|\Gamma_{n}|$, and it will lead to a waste of orders of the filter (the reason will be discussed in Sec. 3.1.1).

- In [26], the gradient descent method is used to optimize (4) (making (4) approximates the objection reflection function). By using the design procedure based on this method, the dimensions of the filter can be predicted. However, (4) is a non-convex function. Therefore, the optimized results are highly related to the initial iterative values [30], and the optimization procedure is time-consuming problem.

To avoid these shortages, this paper proposes a modified method to design a kind of asymmetry stub-loaded microstrip filters.

![Fig. 1. The equivalent model of microstrip filters proposed in [26].](image)

| Notation | Implication |
|----------|-------------|
| $\Gamma_{\text{total}}$ | total reflection coefficient |
| $\Gamma_{n}$ | the $n$th local reflection coefficient |
| $Y_{0}$ | characteristic admittance of the main transmission line |
| $Y_{\text{im}}^{(n)}$ | the $n$th imaginary of the characteristic admittance |
| $Y_{\text{L}}$ | equivalent admittance of the loaded stubs |
| $Y_{x}$ | equivalent admittance of opened-stub or shorted-stub |
| $N$ | order of filter |
| $f$ | frequency |
| $f_{n}$ | the $n$th sampling point of frequency |
| $L_{(n)}$ | length of the $n$th main transmission line |
| $L_{x}^{(n)}$ | length of the $n$th loaded stub |
| $W$ | width of the microstrip line |
| $\theta_{x}^{(n)}$ | equivalent electrical length of the $n$th loaded stub |
| $|S_{11}|$ | return loss |
| $|S_{11}|^{\text{vector form}}$ | return loss vector form |
| $|S_{21}|$ | insert loss |
| $|S_{21}|^{\text{vector form}}$ | insert loss vector form |
| $F_{\text{goal}}$ | optimization goals |
| $F_{\text{goal}}^{\text{vector form}}$ | optimization goals vector form |
| $M$ | number of Monte Carlo experiments |
| $\varepsilon$ | objection function |
| $\varepsilon_{\text{pass}}$ | objection function in the passband |
| $\varepsilon_{\text{stop}}$ | objection function in the stop-band |
| $\varepsilon_{\text{min}}$ | the minimum value of $\varepsilon$ |
| $x^{(n)}$ | flag bit of the $n$th stub |
| $f_{c}^{\text{(pass)}}$ | 3 dB passband cut-off frequency |
| $f_{c}^{\text{(stop)}}$ | 20 dB stop-band cut-off frequency |

Tab. 1. Notations used in this paper.
2.2 Reflection Theory

Firstly, we want to show the calculation process of $|\Gamma_{\text{total}}|$ of the single-stub impedance transformer loaded with one opened- or shorted-stub. The equivalent model is shown in Fig. 2, where $Y_1$ represents the equivalent admittance of the load, $\theta$ represents the equivalent electrical length of the main transmission line, $\theta_x$ represents the equivalent electrical length of the loaded stub, $\Gamma$ represents the reflection coefficient at the input port, and $Y$ depends on whether the loaded stub is an opened-stub or a shorted-stub.

If the $Y_L$, $Y_1$, and $Y_2$ are set to be equivalent to the same value $Y_0$, which means the main transmission line and the load $L$ are matched with each other, the loaded stub and the main transmission line can be regarded as a parallel combination at point $A$. The equivalent admittances $Y_A$ of the opened-stub and the shorted-stub loaded situations are shown in (5), respectively:

$$
\begin{align*}
Y_A^{(\text{opened})} &= Y_0 + jY_0 \tan \theta_x, \quad \text{while} \quad Y_x = 0, \\
Y_A^{(\text{shorted})} &= Y_0 - jY_0 \cot \theta_x, \quad \text{while} \quad Y_x = \infty.
\end{align*}
$$

(5)

For $\Gamma_1$, if we assume that the characteristic impedances of all transmission lines of the multi-section structure equal to each other, the opened-stub loaded and shorted-stub loaded situations are shown as follows [26]:

$$
\Gamma_1^{(\text{opened})} = -\frac{j}{j + 2 \cot \theta_x}, \quad \text{(6)}
$$

$$
\Gamma_1^{(\text{shorted})} = \frac{j}{2 \tan \theta_x - j}. \quad \text{(7)}
$$

Then the $\Gamma$ can be derived as (8):

$$
\Gamma = \frac{|\Gamma_1| + |\Gamma_2| \exp(-2j\theta)}{1 + |\Gamma_1||\Gamma_2| \exp(-2j\theta)}. \quad \text{(8)}
$$

In [26], we assume that the value of $|\Gamma_1| |\Gamma_2|$ is much smaller than 1, therefore, the denominator in (8) can be ignored. In this paper, we use (8) rather than its approximate form to deal with large local reflected situation and derive the more accurate $|\Gamma_{\text{total}}|$ function.

The structure shown in Fig. 1 can be decomposed into $N$ parts of the structure which is shown in Fig. 2. From $N$th part to the 1st part, we can finally derive the $|\Gamma_{\text{total}}|$ function by calculating $N$ times. The derivation steps of $|\Gamma_{\text{total}}|$ from (8) are shown in (9), where $\Gamma_1^{(i)}, \Gamma_2^{(i)}$ and $\Gamma_3^{(i)}$ represents the corresponding reflection coefficient of the $i$th part.

$$
\Gamma^{(N)} = \frac{|\Gamma_1^{(N)}|}{\Gamma_2^{(N)}}, \quad \text{while} \quad \Gamma_2^{(N)} = 0
$$

$$
\Gamma^{(N-1)} = \frac{|\Gamma_1^{(N-1)}| + |\Gamma_1^{(N)}| \exp(-2\theta)}{1 + |\Gamma_1^{(N-1)}||\Gamma_1^{(N)}| \exp(-2\theta)}
$$

$$
\vdots
$$

$$
\Gamma_{\text{total}} = \Gamma^{(1)} = \frac{|\Gamma_1^{(1)}| + |\Gamma_2^{(1)}| \exp(-2\theta)}{1 + |\Gamma_1^{(1)}||\Gamma_2^{(1)}| \exp(-2\theta)}.
$$

(9)

By the $\Gamma_{\text{total}}$ function, the dimensions of the stubs can be obtained directly, and this will be discussed in Sec. 3.1. To verify the feasibility of this modified calculation method, we build a fifth-order shorted-stub loaded microstrip structure (shown in Fig. 3) in the EM simulation software.

All parts of the main line use the same length $L$ and width $W$, as well as the stubs (length $L_x$ and width $W$). The shorted-stubs is achieved by placing via holes at the end of opened-stubs, and the dielectric substrate is chosen to be Rogers RO4350 (dielectric constant 3.66). The dimensions of this structure are shown in Tab. 2.

The calculation results and simulation results are shown in Fig. 4, where $|S_1|$ represents the return loss and $|S_2|$ represents the insert loss.

| Dimension [mm] | $L$ | $L_x$ | $W$ |
|----------------|-----|-------|-----|
| Tab. 2. Dimensions of the structure shown in Fig. 3. |

![Fig. 2. The equivalent model of single-stub impedance transformer loaded with one opened- or shorted-stub.](image1)

![Fig. 3. Fifth-order shorted-stub loaded microstrip structure.](image2)
Figure 4(a) is the calculation results based on the method in [26] (shown in (4)). We can observe that the calculated result obtained by (4) cannot matching the practical $|\Gamma_{\text{total}}|$ curve according to Fig. 4(a), especially with in the stop-band. Fig. 4(b) illustrates that the modified method in this paper can obtain a quite accurate calculation result. The reason of this improvement is shown just after (8), but to achieve this improvement, the cost is also obvious: the $|\Gamma_{\text{total}}|$ function will be an extremely complex formula (9), and it is difficult to calculate the dimensions of the filter by inverse solution through (9).

The equivalent circuit model of the structure shown in Fig. 3 is discussed. For shorted-stub, the rectangular metal line is modeled as $L_1$ and the coupling between the metal line and the ground introduce the capacitor $C_1$. So every shorted-stub is modeled as the parallel LC resonator. Every part of the main line is modeled as the cascade of $C_0$ and $L_0$. The equivalent circuit is shown in Fig. 5. The values of lumped components are optimized in EM software. The equivalent circuit parameters are shown in Tab. 3.

The frequency responses of the circuit simulation are shown in Fig. 6. Compared with the calculated frequency responses obtained by our proposed method, we can find that the frequency responses of the equivalent circuit simulation result fits well with the curves obtained by (9).

At the end of this subsection, we want to show the measurement results of the optimized filter based on the structure shown in Fig. 3. The fabricated filter is shown in Fig. 7, and the practical dimensions of this filter are shown in Tab. 4. The results of the measured S-parameter are shown in Fig. 8.

As depicted in Fig. 8, this filter has an obvious passband from 11 GHz to 39 GHz (3 dB bandwidth). The return loss in the passband of this filter is lower than $-10$ dB, and the out-of-band rejection declines rapidly. The center frequency is around 25 GHz and the relative bandwidth is 112%, therefore, it is an UWB filter.

### Tab. 3. The equivalent circuit parameters of the filter.

| Parameter | $L_0$(nH) | $C_0$(pF) | $L_1$(nH) | $C_1$(pF) |
|-----------|-----------|-----------|-----------|-----------|
| value     | 0.354     | 0.174     | 0.2       | 0.3       |

| Parameter | $W^{(0)}$ | $W^{(1)}$ | $W^{(2)}$ | $W^{(3)}$ | $W^{(4)}$ |
|-----------|-----------|-----------|-----------|-----------|-----------|
| Dimension | 0.6       | 0.55      | 0.45      | 0.45      | 0.55      |

| Parameter | $W^{(1)}$ | $W^{(2)}$ | $W^{(3)}$ | $W^{(4)}$ |
|-----------|-----------|-----------|-----------|-----------|
| Dimension | 0.4       | 0.7       | 0.6       | 0.7       | 0.4       |

| Parameter | $L$ | $L_{\text{x}}$ |
|-----------|-----|---------------|
| Dimension | 3   | 1.6           |

### Tab. 4. Dimensions of the filter.
The optimization goals are shown in Tab. 5. Here we use the 2-norm difference as our objection function $\varepsilon$:

$$
\varepsilon_{\text{pass}} = \left| |S_{11} - F_{\text{goal}}|\right|_2^2 = \sum \left| |S_{11}(f_n) - F_{\text{goal}}(f_n)|\right|^2, \quad (10)
$$

$$
\varepsilon_{\text{stop}} = \left| |S_{21} - F_{\text{goal}}|\right|_2^2 = \sum \left| |S_{21}(f_n) - F_{\text{goal}}(f_n)|\right|^2, \quad (11)
$$

$$
\varepsilon = \sqrt{\varepsilon_{\text{pass}} + \varepsilon_{\text{stop}}}. \quad (12)
$$

where $\varepsilon_{\text{pass}}$ represents the objection function in the passband, $\varepsilon_{\text{stop}}$ represents the objection function in the stop band, $f_n$ represents the $n$th sampling point of frequency, $S_{11}$ represents the vector form of $|S_{11}(f_n)|$, $S_{21}$ represents the vector form of $|S_{21}(f_n)|$ and $F_{\text{goal}}$ represents the vector form of $F_{\text{goal}}$.

For each calculation, the electrical lengths of the stubs are generated randomly, then those values are put into (9) and the difference between the $F_{\text{goal}}$ and the calculated curves obtained by (9) are compared. By executing the Monte Carlo experiments for $M$ times, the minimum value of $\varepsilon$ ($\varepsilon_{\text{min}}$) and its corresponding dimensions of the stubs (the optimal electrical lengths) can be found. By these electrical lengths, the practical lengths of all stubs can be obtained simply. For the BPF being designed here, the relationship between $M$ and $\varepsilon_{\text{min}}$ is shown in Fig. 9.

As shown in Fig. 9, $\varepsilon_{\text{min}}$ decreases with the increase of $M$, and $\varepsilon_{\text{min}}$ is fluctuating while $3000 < M < 15000$, but it is quite steady when $M$ is larger than 15000. Therefore, we choose 15000 as the value of $M$ in the rest of this paper.

So far, the whole calculation procedure has been shown. The flowchart of this procedure is shown in Fig. 10, where $L^{(n)}$ represents the $n$th main line in Fig. 5 (from the left to the right). The value ranges of $N$, $L^{(n)}$ and $L_{\text{t}}^{(n)}$ are shown in Tab. 5. By using this procedure and the design goals in Tab. 5, here a bandpass filter is designed. The final $\varepsilon_{\text{min}}$ is 0.8. Although the error still exists, it is small enough to guide us to design the filter according the later results.
The comparisons between measured and calculated S-parameters are shown in Fig. 11. As it is depicted, the range of $|S_{11}|$ under $-10$ dB is 7.1 GHz–9.1 GHz. The ranges of $|S_{21}|$ under $-20$ dB are 5.5 GHz–6.5 GHz and 10 GHz–11 GHz. The calculated S-parameter results do not perfectly achieve the design goals, but it is enough to guide us in design process. Figure 11 also shows that the calculation results match the measurement results well. The calculated and practical dimensions of this BPF are shown in Tab. 7.

For this BPF, $L_1^{(2)}$, $L_4^{(4)}$ and $L_5^{(5)}$ are chosen as shorted-stubs, and the other stubs are opened-stubs. Whether the stub is chosen to be opened- or shorted-stub can be achieved by setting a flag bit $x^{(n)}$ and we randomly generate it in the design procedure. For example, we can use this $x^{(n)}$ to create a binary detector: if $x^{(n)}$ equals to 1, set the $L_x^{(n)}$ as an opened-stub; if $x^{(n)}$ equals to 0, set the $L_x^{(n)}$ as a shorted-stub.

As shown in Tab. 7, the differences between calculated and practical dimensions are very small, and it means that using our proposed method can save plenty of time while we optimize the filter in the EM software. The stub-loaded structure can be designed to a small size by bending its stubs and this is also a big advantage of our proposed microstrip structure. The fabricated bandpass filter is shown in Fig. 12. As the Fig. 12 shows, by bending the stubs, the space on the substrate board can be saved, and the size of the filter can be shaped into a small size.

![Fig. 11. S-parameter curves of the BPF.](image)

![Fig. 12. The fabricated BPF.](image)

| Parameter | Calculated Dimension [mm] | Practical Dimension [mm] |
|-----------|---------------------------|--------------------------|
| $L_1$     | 5                         | 5                        |
| $L_2$     | 6                         | 6                        |
| $L_3$     | 0.5                       | 0.5                      |
| $L_4$     | 3.6                       | 3.6                      |
| $L_x^{(1)}$ | 11.8                     | 11.6                     |
| $L_x^{(2)}$ | 10.3                     | 10                       |
| $L_x^{(3)}$ | 8.9                      | 8.9                      |
| $L_x^{(4)}$ | 17.6                     | 17.9                     |
| $L_x^{(5)}$ | 16.6                     | 16.2                     |
| $W$       | --                        | 0.6                      |

Tab. 7. Calculated and practical dimensions of the bandpass filter.

![Fig. 9. Relationship between $M$ and $\varepsilon_{\text{min}}$.](image)

![Fig. 10. Flowchart of the calculation procedure.](image)

Tab. 6. Value ranges of $N, L_x^{(n)}$ and $L_x^{(n)}$. 

![Parameter | Calculated Dimension [mm] | Practical Dimension [mm] ](image)
3.2 Discussion

3.2.1 Comparison with Other Method

Some properties and advantages of our proposed method are shown in this subsection.

Compared with the method in [26], our method can calculate more accurately. Furthermore, the method in this paper is less time-consuming compared with the method shown in [26]. The relationships among the number of iterations $M$, $\epsilon_{\text{min}}$, and time cost are shown in Tab. 8. [26] use the logarithmic difference between (4) and goal function as the value of $\epsilon_{\text{min}}$, and we transform the $\epsilon_{\text{min}}$ into the same logarithmic form in Fig. 9. The calculations of both methods are done on the same computer.

Table 8 shows that our modified method can save plenty of calculation time compared with the method in [26]. The reason for this performance is that the gradient descent algorithm is used as optimal method in [26] and the computer need to calculate the partial derivative about (4) for dozens of times in each time of calculation. But for the method proposed in this paper, some further complex calculations are avoided, and we only need to calculate the difference between $|T|_{\text{total}}$ and the goal function each time. For the $\epsilon_{\text{min}}$, the method based on the Monte Carlo experiments converges more quickly than the method based on the gradient descent algorithm. Therefore, the modified method proposed in this paper is more efficient than the method proposed in [26].

| Number of Iterations | Ref. [26] | This Paper |
|----------------------|-----------|------------|
|                      | $\epsilon_{\text{min}}$ | time cost | $\epsilon_{\text{min}}$ | time cost |
| 100                  | 13.2      | 5 s        | 18.3         | 0.6 s     |
| 4000                 | 8.4       | 4 min      | 6.8          | 9.5 s     |
| 8000                 | 7.6       | 9 min10 s  | 6.9          | 51.4 s    |
| 12000                | 7.2       | 14 min10 s | 4.1          | 1 min 34 s|
| 16000                | 5.7       | 21 min30 s | 3.3          | 2 min 10 s|
| 20000                | 4.7       | 29 min20 s | 3.2          | 3 min 6 s |

Tab. 8. Relationship among the number of iterations, $\epsilon_{\text{min}}$ and time cost.

|                      | $f_c^{(\text{pass})}$ [GHz] | $f_c^{(\text{stop})}$ [GHz] | Length [mm] |
|----------------------|-----------------------------|-----------------------------|-------------|
| Chebyshev BPF        | 6.7–9.9                     | 5.5–6.2, 10.2–11            | 27          |
| Proposed BPF         | 7–9.5                       | 5.5–6.5, 10–11              | 2.4         |

Tab. 9. Comparison between Chebyshev BPF and the proposed BPF.

As it is shown in this paper, a filter design method with a high-speed design process is proposed. Here we use the Chebyshev method proposed in [30] to redesign the BPF shown in Sec. 3.1 and compare its results with the results obtained by our proposed method. The detail of the design procedure is explained in [30], and we only give out the results. The structure of the filter based on Chebyshev method is shown in Fig. 13, and the frequency responses of this filter are shown in Fig. 14. As it is depicted in the Fig. 14, by using our method, we can design a filter which has the frequency responses as good as the Chebyshev filter has. The comparison between BPFs obtained by Chebyshev and our proposed methods is shown in Tab. 9, where $f_c^{(\text{pass})}$ represents the 3dB passband cut-off frequency, and $f_c^{(\text{stop})}$ represents the 20dB stop-band cut-off frequency.

But if we focus on the design procedure of the Chebyshev filter, we can find it is much more complex than our proposed method: Firstly, the order of equal-ripple low-pass filter prototype need to be selected while using Chebyshev method. Secondly, the capacitor values and the inductor values of the circuit mode need to be calculated according to the Chebyshev polynomials. Thirdly, the circuit of the LPF prototype need to be translated into a BPF circuit using the method proposed in [30]. Finally, the circuit mode need to be translated into microstrip line using the theory of Richards’ transformation. As a comparison, by using the method proposed in this paper, the dimensions of the filter can be selected directly and automatically.

3.2.2 Dimensions of the Filter

As it is shown in Fig. 12, the structure of our proposed filter is an asymmetric structure. Compared with the symmetric structure proposed in [26], the asymmetric structure can prevent the waste of orders of the filter, and it is effective to miniaturize the shape of the filter by using the asymmetric structure. Let us consider a point symmetrical structure of the proposed BPF (shown in Fig. 15), where the order $N$ of the BPF equals to 10. The $(N/2 – i)$th sub and the $(N/2 + i + 1)$th stubs use the same type of loaded-stub, and their lengths equal to each other $(i = 0, 1, \ldots, N/2 – 1)$. The symmetrical BPF also use the practical dimensions shown in Tab. 7.
The frequency responses of the symmetric structure are simulated, and they are shown in Fig. 16. Compared with the asymmetric structure, the 10-order symmetric BPF can be regarded as a series connection of two asymmetric BPFs. Therefore, the out-of-band rejection of the symmetric BPF can be better than the asymmetric BPF. But the passband can be narrower, and the insert loss in the passband can be worse.

As a conclusion, if the requirements of the filter performance are not extremely high, we can use the asymmetric structure to control the size of the filter.

Here we design a LPF based on the proposed method and compare this LPF with the filters designed in [24] and [29]. The prototype of our LPF is shown in Fig. 17 and its frequency response is shown in Fig. 18. All the parallel stubs we used are opened stubs. The calculated and practical dimensions of this filter are shown in Tab. 10. The comparisons of the 3 dB passband cut-off frequency \( f_{c}^{(\text{pass})} \), 20 dB stop-band cut-off frequency \( f_{c}^{(\text{stop})} \) and the lengths of filters are listed in Tab. 11.

![Fig. 15. Point symmetrical structure of the proposed BPF.](image1)

![Fig. 16. Frequency responses of the symmetrical structure and the asymmetric structure of the BPF.](image2)

![Fig. 17. The fabricated LPF.](image3)

![Fig. 18. S-parameter curves of the LPF.](image4)

| Parameter | Calculated Dimension [mm] | Practical Dimension [mm] |
|-----------|---------------------------|--------------------------|
| \( L_1 \) | 3.9                       | 3.3                      |
| \( L_2 \) | 4.7                       | 3.6                      |
| \( L_3 \) | 4.9                       | 4.9                      |
| \( L_4 \) | 9.7                       | 9.7                      |
| \( L_x^{(1)} \) | 8.6                      | 8.6                      |
| \( L_x^{(2)} \) | 1                        | 1                        |
| \( L_x^{(3)} \) | 6.3                      | 6.6                      |
| \( L_x^{(4)} \) | 3.1                      | 3.1                      |
| \( L_x^{(5)} \) | 0.6                      | 0.6                      |
| \( W \) | –                         | 0.6                      |

**Tab. 10.** Calculated and practical dimensions of our low-pass filter.

| Characteristic | Filter in [28] | Filter in [29] | Filter in this Paper |
|----------------|----------------|----------------|---------------------|
| \( f_{c}^{(\text{pass})} \) [GHz] | 0–4 | 0–3.5 | 0–4 |
| \( f_{c}^{(\text{stop})} \) [GHz] | 6 | 5 | 5 |
| Length [mm] | 14.3 | 2.4 | 2 |

**Tab. 11.** Comparison among low-pass filters in [24], [29] and this paper.
The filter designed in [24] is a stepped impedance filter based on the traditional Butterworth approximation method. The structure is simple but the length is around 7 times longer than our proposed filter, and its stop-band attenuation decays slowly. The filter shown in [29] is designed based on spiral compact microstrip resonant cells (NSCMRCs), and this filter has high performance in the stop-band and can be fabricated in small size. The structure of this filter is quite complex, so it is difficult to design and fabricate. Our filter eliminates the defects and contains the advantages mentioned above. Besides, our filter can perform well both in the passband and the stop-band and the size of our filter can be very small.

In [26], a dual-band filter with narrow passbands has been designed by using the reflection theory. And simulated result shows that the proposed method in [26] is effective under narrow band and multi-band situation. The dual-band filter proposed in [26] is shown in Fig. 19.

As a modified version of the method proposed in [26], the method proposed in our paper also can be used to design narrow band and multi-band filters. The parameters of the dual-band filter proposed in [26] are substituted into (9), and the frequency responses of this filter are calculated.

4. Conclusions

In this paper, a novel method to design microstrip filters is proposed, and the method is based on the reflection theory and the calculation procedure of filter dimensions is based on the Monte Carlo experiments. The derivation procedure of $|\Gamma_{total}|$ is shown in Sec. 2 and the details of design procedure are shown in Sec. 3. Three filters are designed and fabricated. In Sec. 2.2, an UWB filter is designed to show that the calculation of $|\Gamma_{total}|$ is accurate; A bandpass filter is designed in Sec. 3.1 to show that our reverse calculating procedure is feasible; A low-pass filter is designed and compared with some prior arts in Sec. 3.2 to show the advantages of our proposed method and filter. All three above filters perform well, and the calculated and practical dimensions of stubs show a satisfactory fitting degree. At the end of Sec. 3, the frequency responses of a dual-band filter shows that our proposed method not only apply to the wide-band situation but also to the narrow band and multi-band situations. Besides, the design procedure of our method is simple and can calculate quickly compared with the method in [26]. The structure of designed filters is simple and their size is small by using our method. In conclusion, our method can be a good guidance in all kinds of microstrip filters design process.

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