Quartz-Tuning-Fork-Enhanced Spectroscopy Based on Fast Fourier Transform Algorithm

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In this paper, a gas sensing technique based on quartz-crystal-tuning-fork-enhanced spectroscopy (QCTFES) and wavelength modulation spectroscopy (WMS) is reported. To explore the capabilities of this technique, a near-infrared (NIR) diode laser emitting at 1,653 nm and a QCTF-based photoelectric detector are developed for measuring trace methane (CH4). For signal processing, a fast and effective signal analysis method based on the fast Fourier transform (FFT) algorithm is proposed for extracting the absorption intensity signal of the QCTFES-WMS, instead of a lock-in amplifier used for harmonic signal demodulation in traditional QCTF-based detection techniques. Primary laboratory results indicate that an excellent linearity response of CH4 concentration and optical power levels are founded, and a detection limit of 64 ppm is achieved with a 1-s averaging time, which can be further improved to 9 ppm at an optimal integral time of 250 s. Improvements in sensitivity and detectivity can be significantly achieved by using laser sources with higher output power. Compared to traditional WMS technique-based semiconductor photodetectors, the room-temperature QCTF-based WMS shows significant advantages of super-broadband wavelength response, much cheap and tiny.

Keywords: laser spectroscopy, QCTFES, FFT, WMS, Trace gas detection

INTRODUCTION

Global warming and climate change are two current key issues to all countries in the world. To better understand the sources and sinks of atmospheric greenhouse gases and other trace gases, real-time and in-site monitoring of their atmospheric concentration change is essential in developing atmospheric models for a scientific understanding of atmospheric chemistry and environmental change [1]. Atmospheric methane (CH4) is a significant greenhouse gas by influencing tropospheric ozone (O3) and stratospheric water vapor (H2O), with a warming impact of 25 times greater than carbon dioxide (CO2), although its average mixing ratio ∼1.8 parts per million by volume (ppmv) in earth atmosphere is much lower than CO2 (∼460 ppm). Moreover, CH4 is also well-known as an industrial safety hazard, especially in the natural gas and coal mining industry, and a biomarker linked to gut bacteria in clinical medicine [2]. Therefore, various analytical techniques have been developed for methane detection, including chemical method and optical techniques. In comparison, optical detection techniques provide the features of non-invasive, point-of-care quantification of gas concentration without any sample preparation. The main laser-based spectroscopic detection techniques, such as multi-path cell-based tunable
diode laser absorption spectroscopy (MPC-TDLAS), high-precision optical cavity spectroscopy [i.e., cavity ring down spectroscopy (CRDS) or cavity-enhanced absorption spectroscopy (CEAS)], integrated cavity output spectroscopy (ICOS), photoacoustic spectroscopy (PAS) or quartz-enhanced photoacoustic spectroscopy (QEPAS), intracavity absorption spectroscopy, and optical frequency comb cavity-enhanced absorption spectroscopy, provide the significant merits of high sensitivity and high selectivity [3].

The spherical CH₄ molecule exhibits very complex infrared spectrum features with numerous perturbations to energy levels and band interactions. Four candidate regions near the mid-infrared (MIR) atmospheric windows of ~3.3 µm (ν₃ band) and ~7.7 µm (ν₁ band) and two near-infrared (NIR) regions of ~2.3 µm and ~1.65 µm are suitable for methane detection, as shown in Figure 1; the simulated data are taken from the high-resolution atmospheric transmission molecular absorption (HITRAN16) database [4]. The fundamental absorption line transitions in the MIR spectral regions are stronger by two orders of magnitude than those of the first overtone band in the NIR regions. Sensitive CH₄ detection has been reported using quantum cascade lasers (QCLs) and interband cascade lasers (ICLs) [5–9]. However, the MIR laser sources and optoelectronic detectors are still costly compared to the devices operating in NIR, typically an infrared mercury cadmium telluride (MCT) detector (VIGO Systems) [10]. Due to the essential factors of semiconductor materials, their wavelength response range is limited, and the detectivity is the inverse proportion to the bandwidth. For this issue, QCTF-based photoelectric detection technology was proposed for mid-infrared gas absorption spectroscopy [11–13]. Unlike the common QEPAS technique (QCTFs as optoacoustic transducers) [14], this technique exploits quartz piezoelectric properties of QCTFs to realize the transformation of the photoelectric signal, which makes the cost-efficient detector suitable for the whole electromagnetic radiation range. Recently, it has been proved to have a photothermal effect due to the photo-thermo-elastic conversion process [15] and was demonstrated with higher sensing capability than traditional TDLAS and QEPAS [16, 17]. Moreover, the QCTF-based detection technique has been successfully demonstrated with significant advantages for standoff photoacoustic spectroscopy [18, 19]. QCTF-based detection techniques are especially attractive for developing portable spectroscopic gas sensors. However, a commercial lock-in amplifier (typical Stanford Model SRS-830 or Zurich Instruments model MFLI) was usually used to extract the amplitude information from the tuning fork signal in

![Figure 1](hitran-based_simulation_of_methane_and_other_species_absorption_lines.png)
all the related works [15–19], which obviously limits a need for developing miniaturized and inexpensive sensors.

In this work, we report a high-sensitivity gas sensing technique based on a DFB diode laser emitting around 1.653 µm and a QCTF-based wavelength modulation spectroscopy (WMS). Unlike the common QEPAS technique, where the laser beam is passed through the gap between the tines of QCTF, in this work, a standard QCTF used a photoelectric detector by utilizing its piezoelectric effect and resonance effect to measure the light intensity signal. Unlike a common WMS technique, a lock-in amplifier is used for harmonic signal demodulation [20]; here, the absorption intensity signal of the QCTF detector is directly extracted by Fast Fourier transform (FFT) algorithm analysis. To demonstrate the applicability for spectroscopic application, the reported QCTF-based WMS technique is employed for trace methane gas detection.

**EXPERIMENTAL DETAILS**

The experimental platform is designed as shown in Figure 2. Methane (CH₄) was selected as the target of interest due to its important effect in the fields of atmospheric environment. In order to cover the strong absorption line of CH₄ near 6046.95 cm⁻¹, a near-infrared distributed feedback (DFB) diode laser emitting at 1.653 nm was used as the excitation light source. For gas measurement, a quartz gas absorption cell with 45° wedge windows was used, which is specially designed to avoid etalon effect. A CaF₂ mirror was chosen as a gas cell window for the absorption path because of its good transmittance in the whole infrared region and good chemical stability. The gas cell is equipped with a temperature sensor and pressure gauge for recording experimental parameters [21]. Instead of the traditional semiconductor detector, a homemade photoelectric

![FIGURE 2](image1.png)  
**FIGURE 2** | Schematic configuration of the CH₄ sensor system.

![FIGURE 3](image2.png)  
**FIGURE 3** | Flowchart of two signal demodulation algorithms.
detector based on QCTF was used for measuring the laser signal. After being transmitted from the gas cell, the laser beam was directly focused on the surface of the QCTF tine by an optical lens (CaF$_2$, focal length $f = 25$ cm). When the incident light was modulated or pulsed at the resonant frequency of the QCTF, a periodic driving force induces a mechanical resonance of the tuning fork, which in turn yields a piezoelectric current due to the piezoelectric effect of the quartz material. Since the piezoelectric current generated by the QCTF is very weak, a low-noise preamplifier with a feedback resistor of $10^8 \Omega$ was developed to magnify the current signal. For a more detailed description of the QCTF-based photoelectric detector, we refer the reader to [12]. The QCTF detector signal is digitized by a data acquisition card (NI USB-6212) with 16 bits of resolution and finally transferred to a laptop computer and a custom-written LabView-based program for real-time signal processing. The laser wavelength can be coarsely and finely tuned by varying its operating temperatures and injecting currents, respectively, which was controlled by a homemade electronic module. Generally, absorption spectra are obtained by sweeping the laser wavelength over absorption features of CH$_4$ using a low-frequency tuning waveform generated from a LabView function generator. In this study, wavelength modulation spectroscopy (WMS) and a 2nd harmonic detection technique were selected to enhance the detection sensitivity. Therefore, a sinusoidal dither signal was synchronously added to the laser driver to modulate the laser wavelength with the modulation frequency of $f_0/2$, where $f_0$ was the QCTF resonance frequency. Instead of the custom digital lock-in amplifier (DLIA) technique, here, the 2nd harmonic ($2f$) absorption signal was directly extracted by a self-developed algorithm based on fast Fourier transform (FFT), as shown in Figure 3. Details of comparison for both demodulation methods are made with different CH$_4$ samples, and the calculated signal-to-noise ratio (SNR) are summarized in Table 1. The proposed FFT demodulation algorithm shows better performance than the DLIA technique in terms of calculating efficiency and parameter dependencies, and a SNR enhancement factor between 1.6 and 3.0 is obtained. The results indicate that the lower the SNR of the raw spectral signal, the better the SNR improvement factor.

### RESULTS AND DISCUSSION

In the case of WMS-$2f$ detection technique, the amplitude of the 2nd harmonic signal shows significant dependence on the laser modulation amplitude. In theory, the $2f$ signal amplitude is maximized at a modulation index of 2.2, which is defined as the ratio of modulation amplitude and molecular absorption line width, yielding the highest SNR. Therefore, the influence of the laser modulation amplitude on the WMS-$2f$ signal amplitude is first investigated.

| CH$_4$ concentration (%) | FFT | DLIA | Ratio |
|--------------------------|-----|------|-------|
| 0.12                     | 20.05 | 7.18 | 2.79 |
| 0.22                     | 35.15 | 15.60 | 2.25 |
| 0.35                     | 75.10 | 25.70 | 2.92 |
| 0.50                     | 120.15 | 41.09 | 2.92 |
| 0.75                     | 140.52 | 58.88 | 2.39 |
| 1.10                     | 202.76 | 84.25 | 2.41 |
| 1.45                     | 235.16 | 138.66 | 1.70 |
| 1.95                     | 278.28 | 166.03 | 1.68 |

**FIGURE 4** | The relationship of the WMS-$2f$ signal amplitude of CH$_4$ and laser modulation amplitude at one atmospheric pressure.
and shown in Figure 4, which was measured with a sample gas pressure fixed at 1 atm. The result shows that the normalized WMS-2f signal amplitudes increases with the increase of the modulation amplitude and then presents a downward trend. The optimal modulation amplitude was experimentally found to be 40 mV for 1.63% CH₄.

Unlike the calibration-free direction absorption spectroscopy (DAS), the high-sensitivity WMS technique must be calibrated before the actual application to gas detection. For calibrating
the gas sensor system, a series of CH$_4$ samples were prepared by dilution of a basic primary standard with high pure N$_2$. For example, the experimentally measured WMS-2f spectra with different CH$_4$ mixing ratios are presented in Figure 5. For clarity, the WMS-2f signal amplitude at the central position was plotted as a function of gas concentration and analyzed by a linear regression algorithm, which is shown in Figure 6. A regression coefficient R$^2$ of 0.999 was obtained; this result indicates that the gas sensor system has a good linearity, which shows a good agreement with the theoretical expectation.

To evaluate the stability and repeatability of the sensor system, continuous measurement at a constant concentration
TABLE 2 | Detectable molecule in the near infrared spectral range and the estimated sensitivity.

| Species | Wavelength (nm) | Wavenumber (cm⁻¹) | Line intensity | Sensitivity (ppm@1s) |
|---------|-----------------|-------------------|----------------|---------------------|
| C₂H₂    | 1520.086        | 6578.576          | 1.340E-20      | 7                   |
| CH₄     | 1663.722        | 6046.964          | 1.455E-21      | 64                  |
| CO      | 1568.774        | 6374.406          | 2.146E-23      | 4,300               |
| CO₂     | 1432.045        | 6983.019          | 5.980E-23      | 1,500               |
| H₂S     | 1578.292        | 6343.999          | 1.307E-22      | 710                 |
| NH₃     | 1531.882        | 6528.768          | 2.348E-21      | 40                  |
| H₂O     | 1392.253        | 7181.156          | 1.505E-20      | 6                   |
| HF      | 1340.44         | 7460.257          | 4.975E-21      | 18                  |

of 1.63% CH₄ was made and analyzed using Allan’s variance analysis method. The continuous measurement was operated for several hours. For example, the result (as shown in Figure 7) indicates that a detection limit of ~64 ppm is achieved for a 1-s averaging time, which can be further improved to 9 ppm with increasing average time to 250 s.

Similar to the photoelectric effect of semiconductor detectors, the QCTF detector is based on the piezoelectric effect and resonance effect. Linear response characteristics of the detector on incident light intensity also need to be calibrated. For this experiment, the relationship between the QCTF signal amplitude and incident laser power was observed. Figure 8 presents the dependence of the QCTF detector response on the incident laser power with a CH₄ mixing ratio of 1.63%. A good linearity ($R^2 = 0.998$) between the QCTF detector signal amplitude and the incident laser power was found under the limited experimental conditions. From this result, we can conclude that the sensor sensitivity can be significantly improved by using light sources with higher output power. Moreover, the detectivity of the QCTF detector can be further improved by fabricating it in a vacuum cavity, where the Q factor has the maximum response.

CONCLUSION

In this paper, a trace gas sensing technique based on quartz-tuning-fork enhanced wavelength modulation spectroscopy is reported and demonstrated for CH₄ gas measurement. For signal processing, a fast and effective signal analysis method based on the FFT algorithm is developed, instead of a lock-in amplifier used for harmonic signal demodulation in traditional QCTF-based detection techniques. The proposed FFT demodulation algorithm shows better performance than the DLIA technique in terms of calculating efficiency and parameter dependencies, and a SNR enhancement factor between 1.6 and 3.0 is obtained. To demonstrate the applicability for spectroscopic application, the reported sensor technique is evaluated for methane gas detection.

The results indicate that the QCTF-based WMS gas sensor has an excellent linearity response of the CH₄ concentration and the incident laser power. Under limited experimental conditions, a detection limit of 64 ppm is achieved for a 1-s averaging time, which can be further improved to 9 ppm with an averaging time of 250 s.

Improvements in sensitivity and detectivity can be significantly achieved by using light sources with higher output power or enhancing the Q factor of the QCTF in a vacuum environment. Compared to traditional semiconductor photodetectors, the QCTF-based detector shows a significant advantage of ultra-wide spectral response range. Therefore, the gas sensor system can be simply modified for other species measurement by only replacing with new laser sources at other spectral regions, especially for some gas species with near infrared fingerprint absorption features, such as acetylene (C₂H₂@1520 nm), water vapor (H₂O@1392 nm), hydrogen fluoride (HF@1340 nm), ammonia (NH3@1531 nm), oxycarbide (CO@1568 nm and CO₂@1432 nm), and hydrogen sulfide (H₂S@1576 nm), as summarized in Table 2. Optical fiber coupling characteristics of near infrared DFB diode lasers are preferred as spectroscopic radiation sources, and a W-level output power can be easily realized by combining with an erbium-doped fiber amplifier (EDFA) [22–24]. Furthermore, when compared to the costly MCT detector commonly used in mid-infrared laser spectroscopy, the room-temperature QCFT detector is much cheap and tiny [5, 25–28].

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/supplementary material, further inquiries can be directed to the corresponding author/s.

AUTHOR CONTRIBUTIONS

GY was responsible for the experimental setup and writing. LX was responsible for the data processing and picture editing of this paper. HL was responsible for the revision of the paper. JL was responsible for the verification of the overall experiment. All the authors made positive contributions to the work.

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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