Extrapolating microdomain Ca\(^{2+}\) dynamics using BK channels as a Ca\(^{2+}\) sensor

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Ca\(^{2+}\) ions play crucial roles in mediating physiological and pathophysiological processes, yet Ca\(^{2+}\) dynamics local to the Ca\(^{2+}\) source, either from influx via calcium permeable ion channels on plasmic membrane or release from internal Ca\(^{2+}\) stores, is difficult to delineate. Large-conductance calcium-activated K\(^+\) (BK-type) channels, abundantly distribute in excitable cells and often localize to the proximity of voltage-gated Ca\(^{2+}\) channels (VGCCs), spatially enabling the coupling of the intracellular Ca\(^{2+}\) signal to the channel gating to regulate membrane excitability and spike firing patterns. Here we utilized the sensitivity and dynamic range of BK to explore non-uniform Ca\(^{2+}\) local transients in the microdomain of VGCCs. Accordingly, we applied flash photolysis of caged Ca\(^{2+}\) to activate BK channels and determine their intrinsic sensitivity to Ca\(^{2+}\). We found that uncaging Ca\(^{2+}\) activated biphasic BK currents with fast and slow components (time constants being \(\tau_f \approx 0.2\) ms and \(\tau_s \approx 10\) ms), which can be accounted for by biphasic Ca\(^{2+}\) transients following light photolysis. We estimated the Ca\(^{2+}\)-binding rate constant \(k_b\) (\(\approx 1.8 \times 10^8\) M\(^{-1}\)s\(^{-1}\)) for mSlo1 and further developed a model in which BK channels act as a calcium sensor capable of quantitatively predicting local microdomain Ca\(^{2+}\) transients in the vicinity of VGCCs during action potentials.

Ca\(^{2+}\) ions are perhaps the most important second messenger for triggering a variety of biological functions and often signal in a highly compartmentalized manner. For example, fast release of synaptic vesicles from presynaptic terminals is triggered by Ca\(^{2+}\) influx via voltage-gated Ca\(^{2+}\) channels during action potentials, and the local high concentration of Ca\(^{2+}\) transients at the active zones is critical for gating synchronized fusion of synaptic vesicles (SV) and transmitter release. Experimental measurements of global Ca\(^{2+}\) concentrations with Ca\(^{2+}\) indicators and mathematical modeling extrapolate that the peak Ca\(^{2+}\) concentration seen by the Ca\(^{2+}\) sensor on SVs briefly reach tens and even hundreds of micromole levels. Although tremendous progress in recent years has been made with developing novel fluorescent chemical or protein Ca\(^{2+}\) sensors, limited spatiotemporal resolution of these sensors presents a major challenge to directly read out local Ca\(^{2+}\) transients in real-time.

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Large-conductance Ca-activated potassium channels (BK channels), uniquely sensitive to both membrane potential and intracellular Ca\(^{2+}\), abundantly distributed in the excitable cells, regulate the membrane excitability and electrical signals in response to the Ca\(^{2+}\)-influx from the Ca\(^{2+}\)-permeable channels\(^3,5\). The BK channel encoded by Slo1 gene contains two calcium binding sites in the regulator of conductance for K\(^{+}\) (RCK) domains of the carboxy-terminal region\(^3,6\) and may potentially serve as an ideal sensor of local Ca\(^{2+}\) rise. However, the affinity of these binding sites is primarily determined under the circumstance of Ca\(^{2+}\) uniformly sojourning to its binding sites at equilibrium with very little consideration of dynamics of Ca\(^{2+}\) influx or release. Although elaborate Markov models containing multiple parallel open and closed states have been developed to simulate both voltage- and Ca\(^{2+}\) dependent gating kinetics of BK channels well\(^7,8\), the forward binding rate constant of Ca\(^{2+}\) (k\(_{f}\)) remains unknown, making model parameters too unconstrained to meaningfully profile local Ca\(^{2+}\) dynamics.

Previous experiments in inside-out patch configuration have attempted to directly measure k\(_{i}\) by ultrafast Ca\(^{2+}\) concentration jumps via a piezoelectric stepper of two barrel theta pipette\(^9\), which enables a solution exchange in less than 1 ms. However, the patch membrane usually invaginates into the pipette tip and forms Ω-shape geometry, slowing the diffusion of Ca\(^{2+}\) (~10 ms) to reach the inner face of the membrane patch where the RCK domain of BK channels situates\(^8\).

To extrapolate the local Ca\(^{2+}\) dynamics using BK channels as a sensor, it is therefore necessary to develop a superfast approach of Ca\(^{2+}\) delivery mimicking calcium influxes via calcium channels induced by action potentials, and precisely measure k\(_{i}\) in order to quantitatively describe the kinetics of BK channels to such fast Ca\(^{2+}\) transients. In this study, we have applied laser flash photolysis technique of the caged-Ca\(^{2+}\) compound (e.g. NP-EGTA) to achieve instantaneous Ca\(^{2+}\) rises, which has been widely used for studying Ca\(^{2+}\)-dependent processes such as the secretory responses\(^10\). After a UV flash-induced photolysis, the intracellular Ca\(^{2+}\) concentration have two phases of rise, a fast transient Ca\(^{2+}\) rise with peak concentrations up to tens of micromole from the basal [Ca\(^{2+}\)]\(_{i}\), of ~10–200 nM in sub-milliseconds and a slow uniformly steady-state elevation of global [Ca\(^{2+}\)]\(^11,13\). We took advantage of biphasic properties with laser photolysis of the caged-Ca\(^{2+}\) compound to examine both voltage- and calcium-dependent gating behavior, and determined the Ca\(^{2+}\) forward binding rate k\(_{f}\) for BK. Our results demonstrate that BK channels have higher calcium-sensitivity capable to follow up to tens of μM transient Ca\(^{2+}\) changes 0.1–0.2 ms, and established a quantitative model for its utility as the fast local Ca\(^{2+}\) sensor to profile the local Ca\(^{2+}\) transients during action potential firing.

**Results**

**BK-type currents elicited by flash photolysis of caged-calcium showed biphasic activation.** To directly investigate the Ca\(^{2+}\)-sensitivity of BK-type channels, laser UV flashes were used to release Ca\(^{2+}\) from caged compound NP-EGTA (10 mM) and activate the currents of mSlo1 and several mutants (BK-type) at various voltages. Fig. 1A shows macroscopic currents of mSlo1 channels expressed in HEK293 cells, which were evoked by a UV flash. The UV flashes with duration-time of 0.2 ms were delivered after the whole-cell configuration was formed for three minutes to ensure the pipette solution uniformly diffused into the cell. After the mSlo1 current reached steady-state at a given voltage of +30 mV, a 0.2 ms UV pulse (pink line) was excited to photolyze the caged Ca\(^{2+}\) to increase the intracellular [Ca\(^{2+}\)]\(_{i}\) to ~10 μM in less than 1 ms (blue line). The mSlo1 current was further enlarged by the uncaged Ca\(^{2+}\) but exhibited a biphasic activation process, presumably as a result of uncaging that produced an early instantaneous transient followed by a plateau increase of intracellular Ca\(^{2+}\) concentrations in a spatially uniform manner\(^13\). Such a complex response appears to be independent of Ca\(^{2+}\)-release from internal store because preloading cells with thapsigargin (1 μM TG) in recording pipettes for at least 5 min to deplete internal Ca\(^{2+}\) stores before flash photolysis has no effect (sFig. 1A). Strikingly, a hooked current appeared at the end of the fast phase. The boxed trace (right) showed typical biphasic activation with the fast on-time (τ\(_{f-on}\)) and slow time (τ\(_{f-off}\)) constants of ~0.2, 2.4 and 10 ms, respectively. As an approximation of relative distribution of two components at any given flash, we define the fast and slow proportion of the current as R\(_{f}\) = h\(_{f}\)/h and R\(_{s}\) = h\(_{s}\)/h, respectively (Fig. 1A), where h is the total current. For the summation of two components, we have R\(_{f}\) + R\(_{s}\) = 1.

The biphasic activations do not come from different calcium binding sites in RCK domains of BK channels. To examine whether different calcium binding site produces different component, we next made a series of constructs including mutant 5D5N (high-affinity calcium binding site deletion) and the mutant D362A/D367A (low-affinity calcium binding site deletion). Even though both of the 5D5N and D362A/D367A show similar biphasic currents, R\(_{f}\) (5D5N) is much smaller than R\(_{f}\) (D362A/D367A) (Fig. 1B). The triple mutant D362A/D367A/5D5N showed no current elicited by uncaged-Ca\(^{2+}\) after UV flash (Fig. 1B), indicating that BK channels cannot be activated at 30 mV without calcium binding sites. Furthermore, there was no obviously continuous increase in all the BK-type currents at their steady-state stage, suggesting that uncaged Ca\(^{2+}\) did not cause a Ca\(^{2+}\)-dependent recruitment of BK channels from intracellular pool to membrane surface. This was particularly evident for the triple mutant (D362A/D367A/5D5N) in which a lack of the channel responsiveness to uncaged Ca\(^{2+}\) should have not precluded its trafficking presumably derived from Ca\(^{2+}\) dependent fusion of vesicles in HEK293 cells. Figure 1C summarizes the some results of R\(_{f}\) from mSlo1, aforementioned mutants and several others (sFig. 1B,C). The R\(_{f}\) < 1 of all the mSlo1 and mutants, albeit there is a big difference among their
values of $R_f$, indicates that they all respond to uncaged $\text{Ca}^{2+}$ in a biphasic manner. This may imply a lack of direct relationship between the $\text{Ca}^{2+}$ binding sites and biphasic waveform of BK currents, as each of $\text{Ca}^{2+}$ binding sites produces both phases. Additionally, we found the distinct difference in $R_f$ between the mutants L312A and G311I in sFig. 1B, C, i.e., $R_f(\text{L312A}) > R_f(\text{G311I})$, possibly due to their extremely different gating properties.

The fast proportion of the biphasic activation ($R_f$) is voltage-dependent but $\text{Ca}^{2+}$-independent. To determine whether the $R_f$ is voltage dependent, we stimulated the currents of mSlo1, 5D5N, D362A/D367A, D369G and D369G/5D5N by flash at various voltage, measured their values of $R_f$ and then plotted their $R_f$-V curves (Fig. 2A). The shapes of the $R_f$-V curves are similar to their G-V curves. After fitted $R_f$-V curves to Boltzmann equation, we derived their values of $V_{0.5}$ and slope factor (in mV) [mSlo1: 22.5 and 18.7; 5D5N: 50.1 and 14.9; D362A/D367A: −10.0 and 21.3; D369G: −70.6 and 23.0; D369G/5D5N: −45.3 and 23.5; D369A/5D5N: 24.9 and 20.5] and found that their values of $V_{0.5}$ and slope factors increase by degrees according to the following rank order: D369G, D369G/5D5N, D369A/5D5N, 5D5N, basically in parallel to their corresponding G-V curves, implying that the $R_f$ values can mirror differences in the calcium sensitivity (or $\text{Ca}^{2+}$ forward binding rate) of BK-type channels and their mutants in response to instantaneous $\text{Ca}^{2+}$ rise.

In contrast to the strong voltage-dependence, the $\text{Ca}^{2+}$-dependence of $R_f$ is not apparent, based on the measured $[\text{Ca}^{2+}]_i$ during flash experiments (Fig. 2B). Moreover, both the fast and slow activation time...
Determination of a physical calcium binding rate $k_b$. In previous models for BK-type channels in 10-state gating scheme $^5$ (sFig. 3, sTable 1), the apparent Ca$^{2+}$ binding rate $k_b$ was arbitrarily set as 1 in ms$^{-1}$μM$^{-1}$ under an empirical assumption of Ca$^{2+}$ forward binding rate being around 10$^9$ s$^{-1}$M$^{-1}$ for many Ca$^{2+}$ binding proteins including BK channels. Hence, all rate constants that allow optimal fits of experimental data remain physiologically irrelevant, unless Ca$^{2+}$ forward binding rate ($k_b$) is fixed to constrain relative changes in the open ($k_o$) and closed ($k_c$) dissociation equilibrium constants, which hold the $k_o/k_c$ constant. To acquire the physical Ca$^{2+}$ binding rate $k_b$, we only consider a simplest two-state $C \xrightarrow{k_{b}} O$ model of ligand (Ca$^{2+}$) and receptor (BK) to describe the BK currents evoked by the uncaged Ca$^{2+}$ release, which can be seen as an unit model of the MWC model of BK. Here the C is a closed state, and the O is an open state. The forward rate between C and O is Ca$^{2+}$-dependent. Based on the C-O model, we have the calcium binding rate constant $k_o = 1/\tau (\tau = ([Ca^{2+}]_a + k_d))$ in μM$^{-1}$ms$^{-1}$ (eq. (4)), where $k_d$ is the dissociation equilibrium constant, $[Ca^{2+}]_a$ the mean Ca$^{2+}$ concentration and $\tau$ the activation time constant during the fast rising phase. Given all the $\tau$, $[Ca^{2+}]_a$ and $k_o$, we can calculate the calcium binding rate constant $k_b$ of BK-type channels. As a calculated Ca$^{2+}$ spike was always appearing at the initial phase of Ca$^{2+}$ uncaging (sFigs 3 and 8 and $[Ca^{2+}]_a$ calculation in Online Methods), we define the $[Ca^{2+}]_a$ as an averaged $[Ca^{2+}]_a$ over the fast rising time of uncaged Ca$^{2+}$ pulse (Fig. 3A–B3). During the rising period of uncaged Ca$^{2+}$, we calculated the fast rising time constant $\tau$ of BK currents, based on a C-O model (eq. (3) in Online Methods). For instance, we obtained the $\tau = 0.15$ ms at 30 mV by a fit of mSlo1 current to eq. (3) (Fig. 3A), the $[Ca^{2+}]_a = 20.89$ μM, averaged from 0 to a Ca$^{2+}$ peak value $[Ca^{2+}]_p = 48.30$ μM.

**Figure 2.** The kinetic characteristics of $R_f$. (A) Voltage-dependence of $R_f$ is plotted for mSlo1 and its mutants as indicated. Left, the $R_f$-V curves of mSlo1 (black), D362A/D367A (red), 5DSN (blue), D369G (green) and D369A/5DSN (cyan) were fitted to the Boltzmann equation $R_f = R_{max}/(1 + \exp(V/V_0.5))$, respectively. Here $R_{max}$ is the maximal fast proportion, $V_0.5$ the voltage of the half maximal fast proportion and $s$ the slope. $V_0.5$ and $s$ in mV are 22.5 ± 14.9 and 14.9 ± 1.7 for 5DSN, 24.9 ± 0.3 and 20.5 ± 3.7 for D369G (n ≥ 5), and 24.9 ± 0.3 and 20.5 ± 3.7 for D369A/5DSN (n ≥ 5), respectively. Right, the representative currents of mSlo1 were obtained at +30 mV, in the presence of ~1 μM (black) and ~10 μM (dark cyan), respectively.

- $\tau_f$ and $\tau_c$ of BK currents are neither voltage-dependent nor Ca$^{2+}$-dependent (sFig. 2A, B). All the data suggest that the uncaged Ca$^{2+}$ transients in HEK cells under our experimental conditions contain both fast and slow components, consistent with those described in artificial conditions $in vitro^{11,12}$, although reliable measurement of $[Ca^{2+}]_a$ can only be made during slow plateau phase (Fig. 1A).
(Fig. 3B1–B3). By the same logic, we estimate $\tau$ to be ~0.15–0.30 ms with a $[\text{Ca}^{2+}]_a = \sim 10–20 \mu\text{M}$ for other BK-type currents (Fig. 3C). The values of $k_d$ for mSlo1 and its mutants can be readily obtained from $\text{Ca}^{2+}$ dose-response curves of BK-type channels (Fig. 3D, sTable 2). Based on the Eq. (4), we thus derived the values of $k_b$ at 30 mV as 0.18 ± 0.04 (n = 6) for mSlo1, 0.057 ± 0.003 (n = 5) for 5D5N, 0.014 ± 0.01 (n = 4) for D362AD367A and 0.26 ± 0.02 (n = 5) for D369G, respectively, as indicated in sTable 2 and Fig. 3E. Here, the $k_b$ value of mSlo1 is near to the typical limitation of calcium binding rates $10^8 \text{s}^{-1}\text{M}^{-1}$. With this new $k_b$, we now for the first time able to constrain the BK model with physiologically meaningful parameters (sFig. 7, sTable 3) and yield optimal match between all the fits and the current traces of mSlo1 channels to drive subsequent non-stationary calculation of local Ca$^{2+}$ transients.

The BK channels can be used as a Ca$^{2+}$ sensor to extrapolate the local intracellular Ca$^{2+}$ from VGCCs in the vicinity of BK channels during action potentials. The $k_b$ value of mSlo1 is over $10^8 \text{s}^{-1}\text{M}^{-1}$, indicating that the BK channel is highly sensitive calcium sensor compared to those mediating the fusion vesicles in chromaffin cells or SVs in central synapses$^{12,17}$. Such a fast binding rate for Ca$^{2+}$ to activate the BK channel justifies it as an optimal Ca$^{2+}$ sensor to extrapolate the local intracellular Ca$^{2+}$ from VGCCs in the vicinity of BK channels during action potentials (APs). As described previously, using the kinetic model for fitting BK current activated by uncaged Ca$^{2+}$, we can reversly extrapolate the local Ca$^{2+}$ profile. To obtain a proof of principle, we transfected Cav1.2 channel alone or in combination with mSlo1 and confirmed that both channels were successfully expressed in HEK293 cells (sFig.9). In
cells co-expressed mSlo1 and Cav1.2, we estimated the number of BK channels and Ca\textsuperscript{2+} channels using a voltage protocol consisted of a pseudo-action-potential (pseudo-AP) and a voltage step to 50 mV, was designed to record the Ca\textsuperscript{2+}-activated currents of BK channels as to acquire the channel number, co-expressed with Cav1.2 in HEK293 cells. Left bottom, the current (black) were obtained in whole-cell patch bathed in the normal saline from a HEK293 cell co-expressed mSlo1 with Ca\textsubscript{v}1.2 subunits. The current (green) was recorded after applying 10 μM Paxilline, of which major component is Ca\textsuperscript{2+} current. The pure BK current (red) was subtracted from the blocked current (green). Right top, a ramp voltage protocol was designed to determine the reversal potential of BK channels in the same case. Right bottom, the current was obtained with no Paxilline (black), with 10 μM Paxilline (green) and the pure BK current was obtained from their difference (red). The holding potential is set at −60 mV to best mimic the resting potential (~−60 mV co-expressing both channels) and minimize the leak or holding currents for voltage-clamp experiments so as to ease the current subtraction procedure for amplitude measurements. (B) Analysis of BK current variance. The scatter plot was obtained from the current variance versus the mean BK currents recorded by a voltage step to 50 mV in (A). Fitted curve (blue) represents the equation: \(\sigma^2 = iI - I^2/N\) with the parameters \(i = 12.9\) pA at 50 mV and \(N = 1662\), where \(\sigma^2\) represents the variance of BK currents, \(i\) the single-channel current amplitude, \(I\) the mean current and \(N\) the number of BK channels. (C) Estimation of local intracellular [Ca\textsuperscript{2+}]\textsubscript{i}. Fitting the 10-state BK model shown in sFig 3 to the pure BK current evoked by the pseudo-AP to determine the intracellular [Ca\textsuperscript{2+}]\textsubscript{i} (cyan). The black point denotes experimental values of BK currents and the pink line is a fit. Green represents the Ca\textsuperscript{2+} current. Voltage protocol used for calculation was plotted at the top, which is composed of dotted and solid lines, indicating that the initial point for calculation started at −4.75 mV to avoid the noise signal as indicated.

Figure 4. Calculating the local intracellular Ca\textsuperscript{2+} based on BK currents. (A) Left top, the voltage protocol, composed of the waveforms of a pseudo-action-potential (pseudo-AP) and a voltage step to 50 mV, was designed to record the Ca\textsuperscript{2+}-activated currents of BK channels as to acquire the channel number, co-expressed with Cav1.2 in HEK293 cells. Left bottom, the current (black) were obtained in whole-cell patch bathed in the normal saline from a HEK293 cell co-expressed mSlo1 with Ca\textsubscript{v}1.2 subunits. The current (green) was recorded after applying 10 μM Paxilline, of which major component is Ca\textsuperscript{2+} current. The pure BK current (red) was subtracted from the blocked current (green). Right top, a ramp voltage protocol was designed to determine the reversal potential of BK channels in the same case. Right bottom, the current was obtained with no Paxilline (black), with 10 μM Paxilline (green) and the pure BK current was obtained from their difference (red). The holding potential is set at −60 mV to best mimic the resting potential (~−60 mV co-expressing both channels) and minimize the leak or holding currents for voltage-clamp experiments so as to ease the current subtraction procedure for amplitude measurements. (B) Analysis of BK current variance. The scatter plot was obtained from the current variance versus the mean BK currents recorded by a voltage step to 50 mV in (A). Fitted curve (blue) represents the equation: \(\sigma^2 = iI - I^2/N\) with the parameters \(i = 12.9\) pA at 50 mV and \(N = 1662\), where \(\sigma^2\) represents the variance of BK currents, \(i\) the single-channel current amplitude, \(I\) the mean current and \(N\) the number of BK channels. (C) Estimation of local intracellular [Ca\textsuperscript{2+}]\textsubscript{i}. Fitting the 10-state BK model shown in sFig 3 to the pure BK current evoked by the pseudo-AP to determine the intracellular [Ca\textsuperscript{2+}]\textsubscript{i} (cyan). The black point denotes experimental values of BK currents and the pink line is a fit. Green represents the Ca\textsuperscript{2+} current. Voltage protocol used for calculation was plotted at the top, which is composed of dotted and solid lines, indicating that the initial point for calculation started at −4.75 mV to avoid the noise signal as indicated.
These observations also provide important insights into the coupling relationship between VGCCs and BK channels or their distance. Because $I_{\text{Cav1.2}}(0\text{mV}) = -1529\text{pA}$ with a single-channel conductance of $2.4\text{pS}$ (at a physiological external concentration of $2\text{mM \text{[Ca}^2\text{]}_\text{e}}$) and a driven force of $60\text{mV}$ with 100% open probability of CavL channels, we can approximate $N(\text{Cav1.2}) = 10618$ per cell. Given the mean total capacitance of a HEK293 cell about $12\text{pF}$ and the membrane capacitance $C_m = 1\text{pF/cm}^2$ \cite{16}, the cell surface area is $1200\mu\text{m}^2$. In other words, the Cav1.2 density is $8.8\text{Cav1.2/\mu m}^2$. As $N(\text{BK}) = 1662$ (Fig. 4B), the BK density is $1.4\text{BK/\mu m}^2$. Assuming random distribution of both channels in HEK cells, we can proximate that each BK channel is surrounded by $6-8\text{Cav1.2}$. Considering a total of $24\text{pM \text{Ca}^2\text{]}_\text{e}}$ during the single AP firing, we can extrapolate that each BK channel is supported with $3-4\text{pA/mM \text{Ca}^2\text{]}_\text{e}}$ per VGCC. In the case of the intracellular 5\text{mM EGTA}, their coupling distance must be within the micro-domain of about $60-80\text{nm}$, on the basis of established $\text{Ca}^{2+}$ dynamics in cases where native tissues do not express BK or low level of voltage-gated $\text{Ca}^{2+}$ channels have been mapped out, BK channels have been shown to be coupled to many mollusk, crustacean, and mammalian cells in vivo and in vitro, and the coupling is different in different cell types. In this study, we established the frame work for these channels to act as an optimal $\text{Ca}^{2+}$ sensor to track fast $\text{Ca}^{2+}$ transients in real-time. Although many organic and protein $\text{Ca}^{2+}$ sensors have been developed to image $\text{Ca}^{2+}$ dynamics in cells with increasing spatiotemporal resolution, it is not only labor intense but also worrisome that the $\text{Ca}^{2+}$ homostasis is perturbed with addition of exogenous reagents and $\text{Ca}^{2+}$ binding proteins. In contrast, BK channels are native to many cells particularly neurons and synapses, we suggest these channels have a broad utility for detecting and profile $\text{Ca}^{2+}$ transients important for local and compartmentalized signaling.

Unlike other studies on steady-state kinetics of BK, we studied the dynamics of BK channel opening to non-stationary mechanical process. By means of laser flash photolysis techniques in this study, we delineated properties of BK currents such as the biphasic activation, voltage-dependent $R_f(V)$ as well as the voltage- and calcium- independent time constants $\tau$ and $\tau_1$ in response to $\text{Ca}^{2+}$ transients that closely resemble physiological $\text{Ca}^{2+}$ influx through voltage-gated calcium channels or release from internal stores. We determined the calcium forward binding rate $k_b$ for BK channels and advanced previously established 10-state BK model with physiologically relevant rate constants. To our best knowledge, the $k_b$ value ($10^9\text{M}^{-1}\text{s}^{-1}$) for BK channels is among the fastest forward binding rates, comparable to fast $\text{Ca}^{2+}$ buffer BAPTA and calmodulin, but 1 or 2 orders of magnitude faster than slow $\text{Ca}^{2+}$ binding proteins calbindin and calretinin \cite{20,23}. Because fast $\text{Ca}^{2+}$ bindings are directly coupled to BK channel openings, currents can readily serve as a sensitive readouts of local $\text{Ca}^{2+}$ transients, presenting advantages over calmodulin based protein $\text{Ca}^{2+}$ sensors (i.e. GCAMPs) \cite{24}. Indeed, using the BK current activated by the $\text{Ca}^{2+}$ influx of Cav1.2 channels co-expressed in HEK cells, we calculated the time course of local $\text{Ca}^{2+}$ in millisecond resolution during a pseudo-AP and the distance between Cav1.2 and BK. These results provide a solid foundation for further exploration of local $\text{Ca}^{2+}$ transients and downstream coupling targets in native cells.

Similar to the experiments of UV-flash uncaging $\text{Ca}^{2+}$, the $k_b$ of mSlo1, D369G and 5D5N except for D362A/D367A are consistent to the rank order of $R_f$, indicating that the greater the $R_f$, the greater the binding rate $k_b$. The anomalous behavior of $R_f(\text{D362A/D367A})$ may come from the specific gating of mSlo1(\text{D362A/D367A}), but further studies are needed to discern its unique properties. Different rates of $\text{Ca}^{2+}$-dependent activation of these mutants may also expand their utility to profile local $\text{Ca}^{2+}$ dynamics in cases where native tissues do not express BK or low level of voltage-gated $\text{Ca}^{2+}$ channels by cell-specific gene targeting approaches with the aim of minimal perturbation to physiologically functions.

Using the genetic method, we demonstrated that the biphasic currents of BK induced by UV flash is not due to different affinity of $\text{Ca}^{2+}$ binding sites of BK-RCK domains, and instead likely originate from biphasic calcium transients as previously demonstrated \cite{11,12}. Dynamic interactions between the robust endogenous calcium buffers and uncaged $\text{Ca}^{2+}$ rise inside cells may underlie such a process \cite{11}. We also noticed that large fast $\text{Ca}^{2+}$ transients induced relatively small BK hook currents, suggesting that the $\text{Ca}^{2+}$ binding rate $k_b$ to BK is the rate-limiting parameter for follow rapid time course of $\text{Ca}^{2+}$ transients during brief APs. It can be envisaged that new mSlo1 mutants with faster $k_b$ can be developed in future studies to overcome the rate limit and more closely track $\text{Ca}^{2+}$ dynamics in real-time.

Taken together we have developed a novel method to calculate the time course of local $\text{Ca}^{2+}$ by measuring BK currents. The core of this approach is to accurately calculate the number of BK channels and to acquire the pure currents of BK using the specific inhibitors of BK. In most cases, the channel number can be obtained by either variance analysis or steady-state current at a known $\text{Ca}^{2+}$ concentration. In cases where the number and topographic distribution of $\text{Ca}^{2+}$ channels have been mapped out, BK may not only serve as a sensitive sensor of local $\text{Ca}^{2+}$ transients but also coupling distance of their micro- /nano-domains in native cells such as neurons and synapses. The utility of BK channels as $\text{Ca}^{2+}$ sensor in native cells depends critically on the distance between BK channels and $\text{Ca}^{2+}$ channels/stores. The distance measurement can be experimentally determined by using electron microscopy of immune gold particles of different size labeling BK channels and calcium channels in the same preparation. Recent developments in super-resolution microscopy may potentially help determine such distances directly if transgenic knock-in mice with BK and calcium channels tagged with different fluorescent proteins are
created. Alternatively, the distance can be quantitatively calculated by injecting different concentrations of Ca\(^{2+}\) buffers such as EGTA and BAPTA into cells to test the degree of attenuation of BK currents. Although these two buffers have similar equilibrium dissociation constants (Kd), the forward binding constant of EGTA is nearly 100 times slower than BAPTA, and hence not effective in capturing Ca\(^{2+}\) ions from sources that are within nanodomain distance (<100nm) whereas BAPTA will be. On the contrary, if Ca\(^{2+}\) source are relatively far in microdomain distance (>100nm), both EGTA and BAPTA can work equally well. Using linearized buffered Ca\(^{2+}\) diffusion models\(^{19,26}\), one can estimate the mean distance between BK channels and Ca\(^{2+}\) source to enable quantitative profile of brief Ca\(^{2+}\) transients in real time during biological activity.

**Online Methods**

**Cell culture and Transfection.** HEK293 cells were cultured in modified Eagle’s medium (DMEM, Gibco) supplemented with 10% fetal bovine serum (FBS, Gibco) at 37°C incubator with 5% CO\(_2\). The day before transfection, cells were transferred into a 24-well plate and transiently transfected using lipofectamine 2000 (Invitrogen) according to manufacturer’s protocol. Recordings were carried out in 1–2 days after transfection.

**Solutions.** All the flash experiments were performed under the normal saline in bath contained the following (in mM): 140 NaCl, 5 KCl, 1.8 CaCl\(_2\), 2 MgCl\(_2\), 10 HEPES and 10 Glucose with the PH adjusted to 7.4. Pipette solution contained the following (in mM): 120 KCl, 10 NP-EGTA, 8.05 CaCl\(_2\), 2 K\(_2\)ATP, 0.4 Fura-4f and 0.4 Mag-Fura-2 and 10 HEPES with the pH adjusted to 7.2. For the experiments of BK current activated by calcium influx, the extracellular solution contained the following (in mM): 150 NaCl, 4 KCl, 2 CaCl\(_2\), 1 MgCl\(_2\), 10 HEPES (pH 7.38); Pipette solution contained the following (in mM): 140 KCl, 10 NaCl, 5 EGTA, 1 MgCl\(_2\), 10 HEPES (pH 7.36). For the experiments of BK kinetic modeling, intracellular solutions with different free Ca\(^{2+}\) were made by mixing (in mM) 160 MeSO\(_3\)K and 10 HEPES with 5 EGTA (for 0\(\mu\)M Ca\(^{2+}\)); 5 EGTA and 3.25 CaCl\(_2\) (for 1\(\mu\)M Ca\(^{2+}\)); 5 HEDTA and 2.99 CaCl\(_2\) (for 10\(\mu\)M Ca\(^{2+}\)); 0.1 CaCl\(_2\) (for 100\(\mu\)M Ca\(^{2+}\)); 0.3 CaCl\(_2\) (for 300\(\mu\)M Ca\(^{2+}\)), with the pH adjusted to 7.0. Free Ca\(^{2+}\) was estimated by the EGTAETC program (E. McCleskey, Vollum Institute, Portland, OR). Pipette solution contained the following (in mM): 160 MeSO\(_3\)K, 2 MgCl\(_2\), 10 HEPES with the pH adjusted to 7.0. All the chemicals were attained from Sigma except that NP-EGTA, Mag-Fura-2 and Fura-4f were from Invitrogen.

**Electrophysiology.** Patch pipettes pulled from borosilicate glass capillaries with resistance of 2–4 meqohms when filled with pipette solution. Macroscopic currents were recorded 3 minutes after the whole-cell patch formed, and another 3 minutes before the next flash in the same patch. All the flash experiments were performed using the EPC-9 patch-clamp amplifier and corresponding software (HEKA, Germany). Currents were typically digitized at 50kHz and filtered at 8.9 kHz (Bessel) to reduce the impact caused by filter settings\(^{27}\). UV excitation light source (Rapp OptoElectronic, Germany) was used to uncage the intracellular Ca\(^{2+}\) and calcium concentration signals were recorded by measuring the fluorescence ratios of 340/380 nm light provided by monochromatic light source (TILL Photonics, Germany). During the Flash experiments recordings, the laboratory should be kept in a dark environment to prevent light pollution, and both of the inside-out and the whole-cell patch experiments were performed in normal saline solutions. All experiments were performed at room temperature (22–24°C).

**Western Blot.** mSlo1 and C-HA tagged Ca\(_{1.2}\) in pcDNA3.1 were co-expressed in HEK293 cells, 24hrs after transfection, the cells were lysed (lysis buffer contained 20mM Tris-HCl/pH 7.5, 150mM NaCl, 1% NP-40, 0.1% Triton X-100, 0.2mM phenylmethylsulfonyl fluoride and protease inhibitors. After vertical rotated at 4°C for 1h, the lysed cells then were high-speed centrifuged (12, 000 rpm) at 4°C for 30 min. The supernatants were then added loading buffer and boiled at 60°C for 10 min. Proteins in the lysate were separated on polyacrylamide gels and transferred to a nitrocellulose membrane. After blocking with 5% nonfat milk in 0.1% Tween 20 in Tris-buffered saline, the blots were probed with mouse monoclonal anti-Slo1 antibody (Abcam, ab99046) and mouse monoclonal anti-human HA antibody (Millipore, 05-904), respectively. Horseradish peroxidase-coupled goat anti-mouse IgG was used as the secondary antibody for the blots. The membrane was washed with 0.1% Tween 20 in Tris-buffered saline, and proteins were visualized with an enhanced chemiluminescence detection system.

**Data analysis and simulation.** Recording data were analyzed with IGOR (Wavemetrix, Lake Oswego, OR), Clampfit (Axon Instruments, Inc.) and Sigmaplot software (SPSS, Inc.). Unless stated otherwise, the data are presented as mean ± S.D. Calculations of parameters for the kinetic modeling were solved numerically, using a Q-matrix and a particle swarm optimization-golden section search (PSO-GSS) algorithms. The global fitting routines were written and executed with software CeL (Huazhong University of Science and Technology, Wuhan, Hubei, China), compiled with the C++ compiler to run under Windows XP.

The G-V curves of BK-type channels were fitted to the single Boltzmann equation:

\[
G = \frac{G_{max} \cdot [Ca^{2+}] / K_{D}}{1 + [Ca^{2+}] / K_{D}}
\]
where \( V_{st} \) is the voltage at which the conductance \( (G) \) is half the maximum conductance \( (G_{\max}) \) and \( \kappa \) is a factor affecting the steepness of the activations. The equilibrium open probability \( P_o \) can be written as:

\[
P_o = \frac{1}{1 + \exp\left(\frac{V - V_{st}}{\kappa}\right)}.
\]

where \([\text{Ca}^{2+}]\) is the intracellular calcium concentration, \( K_C \) the closed \( \text{Ca}^{2+} \) dissociation constant, \( K_C \) the open \( \text{Ca}^{2+} \) dissociation constant, \( F \) Faraday's constant, \( R \) the universal gas constant, \( T \) temperature, \( V \) voltage, \( Q \) the equivalent gating charge associated with the closed-to-open conformational change, and \( I(0) \) the open-to-closed equilibrium constant in the absence of bound \( \text{Ca}^{2+} \) at 0 mV.

For the \( C_{\text{K}+}[\text{Ca}^{2+}] \) model, we have

\[
I = I_{\text{max}} \cdot (1 - e^{-t/\tau})
\]

\[
k_b = \frac{1}{\tau \cdot ([\text{Ca}^{2+}]_o + k_d)}
\]

where \( I \) denotes the BK current, \( I_{\text{max}} \) the maximum BK current, \( a \) the amplitude of BK current, \( \tau \) the activation time constant during the fast rising phase, \( k_b \) the calcium binding rate constant, \( k_d \) the dissociated equilibrium constant and \( [\text{Ca}^{2+}]_o \) the averaged intracellular \( [\text{Ca}^{2+}] \).

**[Ca^{2+}] calculation.** Kinetic modeling is an important method for better understanding the electrophysiological characteristics of ion channels. In kinetic modeling, the membrane current \( I(t) \) can be depicted as a function of membrane voltage \( V(t) \) and calcium concentration \([\text{Ca}^{2+}]_i(t)\),

\[
I = G(V, [\text{Ca}^{2+}]) \times (V - E),
\]

where \( G(V, [\text{Ca}^{2+}]) \) represents the whole-cell conductance determined by both the voltage and calcium concentration and \( E \) the reversal potential of channels. Unlike the Hodgkin-Huxley model, the total conductance is described as follows:

\[
G(V, [\text{Ca}^{2+}]) = n \times \hat{g} \times P_o(V, [\text{Ca}^{2+}]),
\]

where \( n \) is the number of ion channels on the membrane, \( \hat{g} \) the single-channel conductance, and \( P_o(V, [\text{Ca}^{2+}]) \) the open probability, which can be solved from the kinetic models of channels with numerical method of differential equation (Runge-Kutta or Q-Matrix).

Combining eq. (5) and eq. (6), we can get

\[
I = n \times \hat{g} \times P_o(V, [\text{Ca}^{2+}]) \times (V - E).
\]

In eq. (7), given \( I, n, \hat{g}, V, E \) and \( P_o() \), the instantaneous \([\text{Ca}^{2+}]_i \) can be determined. However, unlike \( I, V, E \), and \( \hat{g}, P_o(V, [\text{Ca}^{2+}]) \), \( n \) is hardly to be obtained directly through existing experimental techniques.

In this paper, a novel three-step method is proposed to obtain the instantaneous \([\text{Ca}^{2+}] \) value. Firstly, \( P_o() \) is determined based on the Monod-Wyman-Changeux (MWC) model. Then, \( n \) is obtained through the steady-state calculation of eq. (7). Finally, the instantaneous \([\text{Ca}^{2+}]_i \) can be obtained through eq. (7). The details of method are executed in three steps as follows (sFig. 8A):

1. \text{Step one: establishing the MWC model for BK-type channels and calculating its open probability } \( P_o(V, [\text{Ca}^{2+}]) \).

Models of BK-type channels are an allosteric MWC model of 10-state C5-O5 with 11-parameters, which reflects the voltage- and \( \text{Ca}^{2+} \)-dependent open probability \( (P_o(V, [\text{Ca}^{2+}])) \). Firstly, a set of activated and deactivated data at different \([\text{Ca}^{2+}] \) level (fits) were fitted to a MWC model to determine the eleven unknown parameters with the software CeL\textsuperscript{15} (sFig. 8A). This model can be used to calculate the channel \( P_o(V, [\text{Ca}^{2+}]) \).

2. \text{Step two: determining the channel number } \( n \) \text{ in eq. (7)}

This step is performed at steady state. From the eq. (7), the steady-state current \( I_{\text{steady}}(\infty) \) can be depicted as

\[
I_{\text{steady}}\equiv n \times \hat{g} \times P_o(V, [\text{Ca}^{2+}]_{\text{steady}}) \times (V - E),
\]
where \( I_{\text{steady}}^+[\text{Ca}^{2+}]_{\text{steady}} \), \( V \) and \( E \) were measured from experiment. \( P_n(V, [\text{Ca}^{2+}]) \) was determined from the model in step 1. Therefore, we got \( n \) from the eq. (8).

3. Step three: Instantaneous \([\text{Ca}^{2+}] \), calculation

After step 1 and 2, the instantaneous \([\text{Ca}^{2+}] \) can be determined based on the eq. (7) with all other known parameters. However, the analytical solution of the instantaneous \([\text{Ca}^{2+}] \) is difficult to obtain because \( P_n(V, [\text{Ca}^{2+}]) \) is in a form of complex differential equation system.

Here, instantaneous \([\text{Ca}^{2+}] \), calculation is treated as a numeric optimization problem, where \([\text{Ca}^{2+}](t)\) is the only parameter to be optimizes. Here the Q-Matrix method is used to calculate the values of \( P_n \) and \( I(t) \) at the time \( t \), and then to minimize the error between the calculated \( I(t) \) and the actual \( I(t) \) by an optimization algorithm of Evolution Strategy (ES).

The ES is a class of numeric optimization techniques based on the ideas of adaption and evolution. A group of floating point solution candidates evolves with search operators, such as selection, recombination and mutation. In common with simple genetic algorithms, the operators are applied in a loop. An iteration of the loop is called a generation. The sequence of generations is continued until a termination criterion is met. The flowchart of ES is shown in sFig. 8B.

When calculation begins, \([\text{Ca}^{2+}](t)\), measured at time 0 in experiment is set as the initial steady-state \([\text{Ca}^{2+}] \). Then \([\text{Ca}^{2+}](t)\), at next time step can be got through ES optimization. This process goes step by step until \([\text{Ca}^{2+}](t)\), at all time step is gotten.

It is obviously that the method in step 2 used to determine the channel number in the UV-flash case completely differs from the variance method\(^{12}\) used in the pseudo-AP case.

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Author Contributions
Z.W., L.Y.W., M.D., J.C. and J.D. conceived and designed the experiments; P.H. and Y.W. performed the flash experiments, H.L., G.Z. and W.Z. performs the action potential experiments; F.X. M.D. and M.Y. designed the method used to calculate constant $k_b$ and local calcium concentration; F.X., M.Y. and W.W. wrote the program used in analysis; P.H., F.X. and H.L. Analyzed the data; J.D. and L.Y.W wrote the paper. All authors reviewed the manuscript.

Additional Information

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