Existence of a lower bound for the distance between point masses of relative equilibria for generalised quasi-homogeneous $n$-body problems and the curved $n$-body problem

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Abstract

We prove that if for relative equilibrium solutions of a generalisation of quasi-homogeneous $n$-body problems the masses and rotation are given, then the minimum distance between the point masses of such a relative equilibrium has a universal lower bound that is not equal to zero. We furthermore prove that the set of such relative equilibria is compact and prove related results for $n$-body problems in spaces of constant Gaussian curvature.

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1 Introduction

By $n$-body problems we mean problems where we are tasked with deducing the dynamics of $n$ point masses. The study of such problems has obvious applications to fields such as celestial mechanics, chemistry, atomic physics and crystallography (see for example [1], [5], [6], [7], [8], [9], [10], [17], [19], [22], [23], [24], [25], [41], [42] and the references therein). The $n$-body problems discussed in this paper are the $n$-body problem in spaces of constant Gaussian curvature and a generalisation of quasi-homogeneous $n$-body problems, which we will call generalised quasi-homogeneous $n$-body problems for short:

Definition 1.1. Let

$$M^k_\sigma = \{(x_1, \ldots, x_{k+1}) \in \mathbb{R}^{k+1} | x_1^2 + \ldots + x_k^2 + \sigma x_{k+1}^2 = \sigma\},$$

where $\sigma$ equals either $+1$, or $-1$ and for $x, y \in M^k_\sigma$ define the inner product

$$x \odot y = x_1 y_1 + \ldots + x_k y_k + \sigma x_{k+1} y_{k+1}.$$

By the $n$-body problem in spaces of constant Gaussian curvature, henceforth referred to as the $n$-body problem in spaces of constant curvature, or curved $n$-body problem (see [22], [23] and [24]), we mean the problem of finding the dynamics of $n$ point particles with respective masses $m_1, \ldots, m_n$ and coordinates $q_1, \ldots, q_n \in M^k_\sigma$, $k \geq 2$, as described by the system of differential equations

$$\ddot{q}_i = \sum_{j=1, j \neq i}^{n} \frac{m_j (q_j - \sigma (q_i \odot q_j) q_i)}{(\sigma - \sigma (q_i \odot q_j)^2)^{\frac{3}{2}}} - \sigma (q_i \odot \dot{q}_i) q_i, \quad i \in \{1, \ldots, n\}. \quad (1.1)$$

Definition 1.2. By generalised quasi-homogeneous $n$-body problems we mean problems where we have to find the dynamics of $n$ point particles $q_1, \ldots, q_n \in \mathbb{R}^{k+1}$, $k \geq 1$ with respective masses $m_1, \ldots, m_n$ as described by the system of differential equations

$$\ddot{q}_i = \sum_{j=1, j \neq i}^{n} m_j (q_j - q_i) f(||q_j - q_i||), \quad i \in \{1, \ldots, n\}, \quad (1.2)$$

where $f : \mathbb{R}_{>0} \to \mathbb{R}$ can be any continuous function with the property that $f(x)$ and $xf'(x)$ are bounded and differentiable for $x$ away from 0 and...
\[ \lim_{x \to 0} f(x) = \pm \infty. \]

If \( f(x) = x^{-\frac{3}{2}} \) and \( k = 2 \), then we speak of the \textit{classical} \( n \)-body problem. If \( f(x) = ax^{-\alpha} + bx^{-\beta} \), \( \alpha, \beta \in \mathbb{R}_{>0}, a, b \in \mathbb{R} \), then we speak of a \textit{quasi-homogeneous} \( n \)-body problem. The reason that \( xf'(x) \) needs to be bounded for \( x \) away from zero is a technical one: In the proof of Theorem 1.4, \( \|Q_{2r} - Q_{jr}\|_{B_{jr}} \) needs to be bounded. \( B_{jr} \) depends on \( f' \) and \( \|Q_{2r} - Q_{jr}\| \) may be unbounded. \( xf'(x) \) being bounded for \( x \) away from zero is sufficient to ensure that \( \|Q_{2r} - Q_{jr}\|_{B_{jr}} \) is bounded and does not impose very strong restrictions on the generality of our problem.

Research into \( n \)-body problems for spaces of constant Gaussian curvature goes back as far as the 1830s, when Bolyai and Lobachevsky (see [4] and [38] respectively) independently proposed a curved 2-body problem in hyperbolic space \( \mathbb{H}^3 \). Since then, \( n \)-body problems in spaces of constant Gaussian curvature have been investigated by mathematicians such as Dirichlet, Schering (see [45], [46]), Killing (see [29], [30], [31]), Liebmann (see [34], [35], [36]) and Kozlov and Harin (see [32]). However, the successful study of \( n \)-body problems in spaces of constant Gaussian curvature for the case that \( n \geq 2 \) began with [22], [23], [24] by Diacu, Pérez-Chavela and Santoprete. After this breakthrough, further results for the \( n \geq 2 \) case were then obtained in [7], [11], [12], [13], [14], [15], [18], [20], [20], [27], [51], [52] and [53]. For a more detailed historical overview, please see [12], [13], [14], [16], [18], or [22].

Quasi-homogeneous \( n \)-body problems for general values of \( n \) started with [10] by Diacu and can be applied to many fields ranging from celestial mechanics to atomic physics to chemistry to crystallography (see [23]). For examples see [5], [6], [8], [9], [10], [17], [19], [33], [37], [39], [40], [41], [42], [43], [54] and the references therein.

Solutions to \( n \)-body problems that constitute point configurations that retain their size and shape over time are called \textit{relative equilibria}. However, following [13], we will use a slightly more general definition: In Euclidean space, any configuration \( q_1(t),...,q_n(t) \) that retains its size and shape over time can be written as \( q_1(t) = T(t)Q_1,\ldots,q_n(t) = T(t)Q_n \), where \( Q_1,\ldots,Q_n \in \mathbb{R}^k \) are constant vectors and \( T(t) \) is a block matrix of rotation matrices. Note that for any rotation matrix \( T(t) \) and \( x, y \in \mathbb{R}^k \), we have that \( \langle T(t)x, T(t)y \rangle = \langle x, y \rangle \), as rotations preserve the distance between points and angles between lines. Consequently, all matrices \( T(t) \) for which \( \langle T(t)x, T(t)y \rangle = \langle x, y \rangle \) are norm preserving (take \( x = y \)) and thus preserve angles (as for any vectors \( x, y \), the norms of \( x \), \( y \) and \( x - y \) are preserved and thus the angles of the triangle with sides of length \( \|x\|, \|y\| \) and \( \|x - y\| \) are preserved). On \( \mathbb{R}^k \), \( T(t) \) being
any block matrix of rotation matrices is equivalent with \( \langle T(t)x, T(t)y \rangle = \langle x, y \rangle \). On \( \mathbb{M}_n^k \), if we look for \( T(t) \) for which \( x \odot_k y = (T(t)x) \odot_k (T(t)y) \), that depend on hyperbolic and parabolic functions exist as well (see [13]).

As we do not explicitly use the many different possible expressions for \( T(t) \), it makes sense to therefore define relative equilibria as follows:

**Definition 1.3.** Consider any solution to (1.1), or (1.2) for which the \( q_1, \ldots, q_n \) can be written as \( q_1(t) = T(t)Q_1, \ldots, q_n(t) = T(t)Q_n \), where \( Q_1, \ldots, Q_n \in \mathbb{R}^{k+1} \) are constant and \( T(t) \) is a time dependent, invertible \( k \times k \) matrix for which there exist constants \( c_1, c_2 \in \mathbb{R}_{>0} \) such that for all \( x \in \mathbb{R}^k \) \( c_1 \|x\| \leq \|T(t)^{-1} \dot{T}(t)x\| \leq c_2 \|x\| \). If for all \( x, y \in \mathbb{R}^k \) we have that \( \langle T(t)x, T(t)y \rangle = \langle x, y \rangle \), we call solutions of this type that solve (1.2) a relative equilibrium of (1.2). If for all \( x, y \in \mathbb{M}_n^k \) we have that \( (T(t)x) \odot (T(t)y) = x \odot y \) and \( \|\dot{T}(t)x\| \) is bounded if \( x \) is bounded, then we call solutions of this type that solve (1.1) a relative equilibrium of (1.1). For generalised quasi-homogeneous \( n \)-body problems, we call the set of all such configurations that are equivalent under rotation and scalar multiplication a class of relative equilibria. It should be remarked that the property that \( c_1 \|x\| \leq \|T(t)^{-1} \dot{T}(t)x\| \leq c_2 \|x\| \) does not follow directly from the fact that \( x \odot_k y = (T(t)x) \odot_k (T(t)y) \). Checking the different possible cases in [13] shows easily that the given properties of \( T(t) \) are true and saves us many pages of writing out block matrices.

Relative equilibria can tell a great deal about the physical space for which their respective \( n \)-body problems have been defined: A prime example of how much information can be deduced by studying relative equilibria, comes from celestial mechanics: It was proven in [22] and [23] that for the \( n \)-body problem in spaces of constant curvature (i.e. spheres or hyperboloids) relative equilibria that are shaped as equilateral triangles have to have equal masses. This means that our solar system, with the Sun, Jupiter and the Trojan asteroids forming approximately an equilateral triangle and relative equilibrium, is likely flat within the respective area, i.e. has zero Gaussian curvature. For further information on the relevance of relative equilibria, see for example [21] and [44]. Of particular importance is the link with the sixth Smale problem (see [48]), which states that for the classical case, if the equilibria are induced by a plane rotation, the number of classes of relative equilibria is finite, if the masses \( m_1, \ldots, m_n \) are given. This problem is still open for \( n > 5 \) and was solved for \( n = 3 \) by A. Wintner (see [54]), \( n = 4 \) by M. Hampton and R. Moeckel (see [28]) and for \( n = 5 \) by A. Albouy and V. Kaloshin, assuming that the 5-tuple of positive masses belongs to a given codimension.
2 subvariety of the mass space (see [3]). As a potential step towards a proof of Smale’s problem, M. Shub showed in [47] that the set of all classes of relative equilibria, provided they have the same set of masses, is compact. Additionally, Shub proved in the same paper that if the rotation inducing the equilibria is given as well, there exists a universal nonzero, minimal distance that the point masses lie apart from each other.

In this paper, as a logical next step after Shub’s work in [47] and to gain further understanding of the geometry of relative equilibria, we prove Shub’s results when using (1.2) instead of the classical \( n \)-body problem and related results for \( n \)-body problems in spaces of constant curvature. Specifically, we prove that

**Theorem 1.4.** Consider the set \( R_{T,m_1,...,m_n} \) of all relative equilibria of (1.2) with rotation matrix \( T(t) \) and masses \( m_1,..., m_n \). Then there exists a constant \( c \in \mathbb{R}_{>0} \) such that for all relative equilibria \( \{T(t)Q_i\}_{i=1}^n \) in the set \( R_{T,m_1,...,m_n} \), we have that \( \|Q_i - Q_j\| > c \) for all \( i, j \in \{1,...,n\}, i \neq j \).

and consequently that

**Corollary 1.5.** Consider the set \( R_{T,m_1,...,m_n} \) of all relative equilibria (1.2) with rotation matrix \( T(t) \) and masses \( m_1,..., m_n \). If \( \lim_{x \to 0} xf(x) = \pm \infty \) and \( xf(x) \) is bounded for \( x \) away from \( x = 0 \), then there exists a \( C \in \mathbb{R}_{>0} \) such that for all relative equilibria \( \{T(t)Q_i\}_{i=1}^n \) in the set \( R_{T,m_1,...,m_n} \), we have that \( \|Q_i\| < C \) for all \( i \in \{1,...,n\} \).

For the \( n \)-body problem in spaces of constant curvature, proving that relative equilibria form a compact set is pointless for the case that \( \sigma = 1 \) (i.e. the problem is defined on the unit sphere) but it does make sense to investigate whether there exists a universal lower bound for the distance between the point masses. To that extent, we will prove that

**Theorem 1.6.** Let \( \sigma = 1, \epsilon > 0 \) and let \( R_{\epsilon,T,m_1,...,m_n} \) be the set of all relative equilibria \( T(t)Q_1,...,T(t)Q_n \) of (1.1) for which \( \langle Q_i, Q_j \rangle > -1 + \epsilon, \) \( i, j \in \{1,...,n\}, i \neq j \). Then for all \( \epsilon > 0 \) there exists a constant \( c_\epsilon > 0 \) such that for any relative equilibrium solution in \( R_{\epsilon,T,m_1,...,m_n} \) of (1.1), \( \|q_i - q_j\|_k > c_\epsilon \) for all \( i, j \in \{1,...,n\}, i \neq j \) if the masses \( m_1,..., m_n \) and rotation \( T(t) \) are given.

For the \( n \)-body problem on a hyperbola, relative equilibria do not form a compact set with respect to the Euclidean norm. For the negative curvature case, we will prove that
Theorem 1.7. Let \( \sigma = -1 \) and let \( R_{T, m_1, \ldots, m_n} \) be the set of all relative equilibria \( T(t)Q_1, \ldots, T(t)Q_n \) of (1.1). Then for any bounded subset \( W \) of \( R_{T, m_1, \ldots, m_n} \) there exists a constant \( C_W > 0 \) such that for any relative equilibrium solution in \( W \) of (1.1), \( \| q_i - q_j \|_k > C_W \) for all \( i, j \in \{1, \ldots, n\} \), \( i \neq j \) if the masses \( m_1, \ldots, m_n \) and rotation \( T(t) \) are given.

Remark 1.8. Theorem 1.6 and Theorem 1.7 were proven for a very specific subclass of relative equilibria in [53]. In this paper, it should be noted that Theorem 1.6 and Theorem 1.7 hold for all types of relative equilibria (positive elliptic relative equilibria, positive elliptic-elliptic relative equilibria, negative elliptic relative equilibria, negative hyperbolic relative equilibria, negative elliptic-hyperbolic relative equilibria and higher dimensional versions thereof) of the \( n \)-body problem in \( S^k \) and \( H^k \) and that their proofs do not rely on specific properties of the matrix \( T \).

Remark 1.9. The restriction that the relative equilibria on the unit sphere lie in \( R_{\epsilon, T, m_1, \ldots, m_n} \) is needed to potentially exclude sequences of relative equilibria in \( \mathbb{S}^k \) that in the limit can show antipodal behaviour, i.e. have sequences of point masses \( \{Q_{ip}\}_{p=1}^\infty, \{Q_{jp}\}_{p=1}^\infty \) for which \( \lim_{p \to \infty} \langle Q_{ip}, Q_{jp} \rangle = -1 \). Proving, or disproving the existence of such sequences could lead to valuable information about the geometry of the \( n \)-body problem on the unit sphere, but unfortunately lies beyond the scope of this paper.

We will now prove Theorem 1.4 in section 2 then prove Corollary 1.5 in section 3 after which we will prove Theorem 1.6 and Theorem 1.7 in section 4.

2 Proof of Theorem 1.4

Proof. We will prove this theorem for the case that \( \lim_{x \to 0} f(x) = +\infty \) and give an argument at the end of the proof how the argument of the proof can also be used if \( \lim_{x \to 0} f(x) = -\infty \).

Assume that Theorem 1.4 is false. Then there exist sequences \( \{Q_{ir}\}_{r=1}^\infty \) and relative equilibria \( q_{ir}(t) = T(t)Q_{ir}, i \in \{1, \ldots, n\} \) for which we may assume, if we renumber the \( Q_{ir} \) in terms of \( i \) and take subsequences if necessary, the following:

1. There exist sequences \( \{Q_{1r}\}_{r=1}^\infty, \ldots, \{Q_{lr}\}_{r=1}^\infty, l \leq n \) such that \( \|Q_{ir} - Q_{jr}\| \) goes to zero for \( r \) going to infinity if \( i, j \in \{1, \ldots, l\}, 2 \leq l \geq n \).
2. $\|Q_{ir} - Q_{jr}\|$ does not go to zero for $r$ going to infinity if $i \in \{1, ..., l\}$ and $j \in \{l + 1, ..., n\}$.

3. $\|Q_{1r} - Q_{2r}\| \geq \|Q_{ir} - Q_{jr}\|$ for all $i, j \in \{1, ..., l\}$, for all $r \in \mathbb{N}$.

Write $T(t)^{-1} \dot{T}(t) = -A$. Then inserting $q_{ir}(t) = T(t)Q_i, i \in \{1, ..., n\}$ into (1.2), using that for any $x \in \mathbb{R}^k \|T(t)x\| = \|x\|$ and multiplying both sides of (1.2) with $T(t)^{-1}$, gives

$$-AQ_{ir} = \sum_{j=1, j \neq i}^n m_j (Q_{jr} - Q_{ir}) f(\|Q_{jr} - Q_{ir}\|)$$

and consequently

$$-AQ_{1r} = \sum_{j=2}^n m_j (Q_{jr} - Q_{1r}) f(\|Q_{jr} - Q_{1r}\|) \quad (2.1)$$

and

$$-AQ_{2r} = \sum_{j=1, j \neq 2}^n m_j (Q_{jr} - Q_{2r}) f(\|Q_{jr} - Q_{2r}\|). \quad (2.2)$$

Subtracting (2.1) from (2.2) gives

$$A(Q_{1r} - Q_{2r}) = (m_1 + m_2)(Q_{1r} - Q_{2r}) f(\|Q_{1r} - Q_{2r}\|)$$

$$+ \sum_{j=3}^l m_j ((Q_{1r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|) + (Q_{jr} - Q_{2r}) f(\|Q_{jr} - Q_{2r}\|))$$

$$+ \sum_{j=l+1}^n m_j ((Q_{1r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|) + (Q_{jr} - Q_{2r}) f(\|Q_{jr} - Q_{2r}\|)) \quad (2.3)$$

Note that for any $j \in \{1, ..., l\}$ the vectors $Q_{1r} - Q_{2r}, Q_{1r} - Q_{jr}$ and $Q_{jr} - Q_{2r}$ either form a triangle with $\|Q_{1r} - Q_{2r}\|$ the length of its longest side, or the three of them align, meaning the angles between them are zero. Consequently, the angle between $Q_{1r} - Q_{2r}$ and $Q_{1r} - Q_{jr}$ and the angle between $Q_{1r} - Q_{2r}$ and $Q_{jr} - Q_{2r}$ is smaller than $\frac{\pi}{2}$ and thus

$$\lim_{r \to \infty} (Q_{1r} - Q_{jr}, Q_{1r} - Q_{2r}) \geq 0 \quad \text{and} \quad \lim_{r \to \infty} (Q_{jr} - Q_{2r}, Q_{1r} - Q_{2r}) \geq 0. \quad (2.4)$$
Also note that

\[
\sum_{j=l+1}^n m_j ((Q_{1r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|) + (Q_{jr} - Q_{2r}) f(\|Q_{jr} - Q_{2r}\|))
\]

\[
= \sum_{j=l+1}^n m_j ((Q_{1r} - Q_{2r} + Q_{2r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|) + (Q_{jr} - Q_{2r}) f(\|Q_{jr} - Q_{2r}\|))
\]

\[
= \sum_{j=l+1}^n m_j ((Q_{1r} - Q_{2r}) f(\|Q_{1r} - Q_{jr}\|)
\]

\[+(Q_{jr} - Q_{2r}) (f(\|Q_{jr} - Q_{2r}\|) - f(\|Q_{1r} - Q_{jr}\|)))
\]  

(2.5)

Let

\[
B_{jr} = \begin{cases} 
    f(\|Q_{1r} - Q_{2r}\|) - f(\|Q_{1r} - Q_{jr}\|), & \text{if } \|Q_{jr} - Q_{2r}\| \neq \|Q_{1r} - Q_{jr}\| \\
    0, & \text{if } \|Q_{jr} - Q_{2r}\| = \|Q_{1r} - Q_{jr}\|
\end{cases}
\]

Taking inner products on both sides of (2.3) with \(\frac{Q_{1r} - Q_{2r}}{\|Q_{1r} - Q_{2r}\|}\), using (2.4), (2.5) and using that for \(r\) large enough and \(i, j \in \{1, \ldots, l\}\) we have that \(f(\|Q_{1r} - Q_{jr}\|) \geq 0\), gives for \(r\) large enough that

\[
\left( \frac{\langle A(Q_{1r} - Q_{2r}), Q_{1r} - Q_{2r} \rangle}{\|Q_{1r} - Q_{2r}\|^2} \right) \geq (m_1 + m_2) f(\|Q_{1r} - Q_{2r}\|) + 0
\]

\[
+ \sum_{j=l+1}^n m_j f(\|Q_{1r} - Q_{jr}\|) + \sum_{j=l+1}^n m_j \langle Q_{jr} - Q_{2r}, \frac{Q_{1r} - Q_{2r}}{\|Q_{1r} - Q_{2r}\|} \rangle B_{jr}.
\]

Note that \(\left| \frac{\langle A(Q_{1r} - Q_{2r}), Q_{1r} - Q_{2r} \rangle}{\|Q_{1r} - Q_{2r}\|^2} \right| \leq \|A(Q_{1r} - Q_{2r})\| \|Q_{1r} - Q_{2r}\| \|Q_{1r} - Q_{2r}\| \leq c_2\), that \(\frac{Q_{1r} - Q_{2r}}{\|Q_{1r} - Q_{2r}\|}\]

is bounded in norm, that for \(\lim_{r \to \infty} Q_{1r} = \lim_{r \to \infty} Q_{2r} = Q_1\) and \(\lim_{r \to \infty} Q_{jr} = Q_j\)

\[
\lim_{r \to \infty} B_{jr} = \begin{cases} 
    \lim_{r \to \infty} f'(\|Q_{1r} - Q_{jr}\|) \frac{\|Q_{jr} - Q_{2r}\| - \|Q_{1r} - Q_{jr}\|}{\|Q_{1r} - Q_{2r}\|} & \text{if } \|Q_{jr} - Q_{2r}\| \neq \|Q_{1r} - Q_{jr}\| \\
    0 & \text{if } \|Q_{jr} - Q_{2r}\| = \|Q_{1r} - Q_{jr}\|
\end{cases}
\]

which is bounded by construction as by the triangle inequality,

\[
\left| \frac{\|Q_{jr} - Q_{2r}\| - \|Q_{1r} - Q_{jr}\|}{\|Q_{1r} - Q_{2r}\|} \right| = \left| \frac{\|Q_{jr} - Q_{2r}\| - \|Q_{1r} - Q_{jr}\|}{\|Q_{1r} - Q_{2r}\|} \right| \leq \frac{\|Q_{jr} - Q_{2r}\|}{\|Q_{1r} - Q_{2r}\|} = 1,
\]
which means that we have that for some $B > 0$

$$\lim_{r \to \infty} \left| \sum_{j=l+1}^{n} m_j f(\|Q_{1r} - Q_{jr}\|) + \sum_{j=l+1}^{n} m_j \langle Q_{jr} - Q_{2r}, \frac{Q_{1r} - Q_{2r}}{\|Q_{1r} - Q_{2r}\|} \rangle B_{jr} \right| < B$$

for all $r \in \mathbb{N}$ and thus that

$$c_2 \geq \lim_{r \to \infty} (m_1 + m_2) f(\|Q_{1r} - Q_{2r}\|) - B = \infty,$$

which is a contradiction. If $\lim_{x \downarrow 0} f(x) = -\infty$, we can define $g = -f$, rewrite everything in terms of $g$ and repeat the proof of our theorem using $g$ instead of $f$. This completes our proof.

\end{proof}

3 Proof of Corollary 1.5

\begin{proof}
Assume the contrary to be true. Then there exist sequences $\{Q_{ir}\}_{r=1}^{\infty}$, $i \in \{1, ..., n\}$ for which $q_i(t) = T(t)Q_{ir}$ define relative equilibrium solutions of (1.2) and for which there has to be at least one sequence $\{Q_{ir}\}_{r=1}^{\infty}$ that is unbounded. Taking subsequences and renumbering the $Q_{ir}$ in terms of $i$ if necessary, we may assume that $\{Q_{1r}\}_{r=1}^{\infty}$ is unbounded. Then by (2.1),

$$A Q_{1r} = \sum_{j=2}^{n} m_j (Q_{1r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|)$$

(3.1)

with $A = -T(t)^{-1} \ddot{T}(t)$. As the left-hand side of (3.1) is unbounded, the right-hand side must be unbounded as well, which means that there must be $j \in \{2, ..., n\}$ for which

$$m_j (Q_{1r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|)$$

is unbounded if we let $r$ go to infinity. But as

$$\|m_j (Q_{1r} - Q_{jr}) f(\|Q_{1r} - Q_{jr}\|)\| = m_j \|Q_{1r} - Q_{jr}\| f(\|Q_{1r} - Q_{jr}\|),$$

that means that $\|Q_{1r} - Q_{jr}\|$ goes to zero for $r$ going to infinity, which is impossible by Theorem 1.4. This completes the proof.

\end{proof}
4 Proof of Theorem 1.6 and Theorem 1.7

Proof. We will prove Theorem 1.6 and Theorem 1.7 using the same argument:
Assume that either of the theorems is incorrect. Then there exist sequences of relative equilibria \( \{ q_{ir}(t) \}_{r=1}^{\infty} = \{ T(t)Q_{ir} \}_{r=1}^{\infty}, \ i \in \{ 1, \ldots, n \} \), for which, renumbering the \( q_{ir} \) in terms of \( i \) if necessary, there exists an \( l \in \{ 1, \ldots, n \} \) such that

1. \( \| Q_{ir} - Q_{jr} \| \) goes to zero for \( r \) going to infinity if \( i, j \in \{ 1, \ldots, l \} \).

2. \( \| Q_{ir} - Q_{jr} \| \) does not go to zero for \( r \) going to infinity if \( i \in \{ 1, \ldots, l \} \) and \( j \in \{ l+1, \ldots, n \} \).

Inserting these solutions into (1.1), using that for \( x, y \in M^k \), \( (T(t)x) \odot (T(t)y) = x \odot y \), subsequently multiplying both sides of the equation from the left with \( T(t)^{-1}T(t)Q_{ir} + \sigma(q_{ir} \odot \dot{q}_{ir})Q_{ir} = B_{\sigma ir} \) gives

\[
B_{\sigma ir} = \sum_{j=1, j \neq i}^{n} \frac{m_j(Q_{jr} - \sigma(Q_{ir} \odot Q_{jr})Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}}. \tag{4.1}
\]

Multiplying both sides of (4.1) with \( m_i \) and summing over \( i \) from 1 to \( l \) then gives

\[
\sum_{i=1}^{l} m_i B_{\sigma ir} = \sum_{i=1}^{l} \sum_{j=1}^{n} \sum_{j \neq i}^{n} \frac{m_i m_j (Q_{jr} - \sigma(Q_{ir} \odot Q_{jr})Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}}. \tag{4.2}
\]

Note that

\[
2 \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{jr} - \sigma(Q_{ir} \odot Q_{jr})Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}}

= \left( \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{jr} - \sigma(Q_{ir} \odot Q_{jr})Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}} + \sum_{j=1}^{l} \sum_{i=1, i \neq j}^{l} \frac{m_j m_i (Q_{jr} - \sigma(Q_{jr} \odot Q_{ir})Q_{jr})}{(\sigma - \sigma(Q_{jr} \odot Q_{ir})^2)^{\frac{3}{2}}} \right),
\]

which, collecting all terms that are multiplied with \( m_i m_j \), can be rewritten
as

\[ 2 \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{jr} - \sigma(Q_{ir} \odot Q_{jr}) Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}} \]

\[ = \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{ir} + Q_{jr})(1 - \sigma(Q_{ir} \odot Q_{jr}))}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}} \]

\[ = \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{ir} + Q_{jr})(1 - \sigma(Q_{ir} \odot Q_{jr}))}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}} \cdot \frac{\sigma(1 + \sigma(Q_{ir} \odot Q_{jr}))}{\sigma(1 + \sigma(Q_{ir} \odot Q_{jr}))}, \]

giving

\[ 2 \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{jr} - \sigma(Q_{ir} \odot Q_{jr}) Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}} \]

\[ = \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{ir} + Q_{jr})(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}(1 + Q_{ir} \odot Q_{jr})} \]

\[ = \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{ir} + Q_{jr})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}(1 + Q_{ir} \odot Q_{jr})}. \quad (4.3) \]

Let

\[ B_{2\sigma r} = \sum_{i=1}^{l} \sum_{j=l+1}^{n} \frac{m_i m_j (Q_{jr} - \sigma(Q_{ir} \odot Q_{jr}) Q_{ir})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}}. \quad (4.4) \]

Inserting (4.3) and (4.4) into (4.2) gives

\[ \sum_{i=1}^{l} m_i B_{\sigma ir} Q_{ir} = B_{2\sigma r} = \frac{1}{2} \sigma \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{ir} + Q_{jr})}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}(1 + Q_{ir} \odot Q_{jr})}, \]

which means that

\[ \left( \sum_{i=1}^{l} m_i B_{\sigma ir} Q_{ir} - B_{2\sigma r} \right) \odot Q_{1r} = \frac{1}{2} \sigma \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_i m_j (Q_{ir} + Q_{jr}) \odot Q_{1r}}{(\sigma - \sigma(Q_{ir} \odot Q_{jr})^2)^{\frac{3}{2}}(1 + Q_{ir} \odot Q_{jr})}. \quad (4.5) \]
By construction, \( \left( \sum_{i=1}^{l} m_i B_{\sigma ir} Q_{ir} - B_{2\sigma r} \right) \circ Q_{1r} \) and \( \sigma (1 + Q_{ir} \circ Q_{jr}) > 0 \) are bounded for \( r \) going to infinity and \( \lim_{r \to \infty} Q_{ir} \circ Q_{jr} = 1 \) for \( i, j \in \{1, ..., l\} \), so by (4.5)

\[
\lim_{r \to \infty} \left( \sum_{i=1}^{l} m_i B_{\sigma ir} Q_{ir} - B_{2\sigma r} \right) \circ Q_{1r} = \frac{1}{2} \sum_{i=1}^{l} \sum_{j=1, j \neq i}^{l} \frac{m_im_j(Q_{ir} + Q_{jr}) \circ Q_{1r}}{(\sigma - \sigma(Q_{ir} \circ Q_{jr})^2) + \frac{1}{2} \sigma(1 + Q_{ir} \circ Q_{jr})} = \infty, \quad (4.6)
\]

while the left-hand side of (4.6) is bounded, which is a contradiction. This completes the proof. 

\[ \square \]

References

[1] R. Abraham, J. Marsden, (1978) Foundations of Mechanics, Addison-Wesley Publishing Co. Reading, Mass.

[2] Some problems on the classical n-body problem, Celest. Mech. Dyn. Astr. 133, (2012) 369–375.

[3] A. Albouy, V. Kaloshin, Finiteness of central configurations of five bodies in the plane, Annals of Mathematics 176, (2012) 535–588.

[4] W. Bolyai and J. Bolyai, Geometrische Untersuchungen, Hrsg. P. Stäckel, Teubner, Leipzig-Berlin, 1913.

[5] S. Craig, F. Diacu, E.A. Lacomba and E. Perez-Chavela, On the anisotropic Manev problem, J. Math. Phys. 40, (1999) 1–17.

[6] M. Corbera, J. Llibre and E. Pérez-Chavela, Equilibrium points and central configurations for the Lennard-Jones 2- and 3-body problems, Celestial Mechanics and Dynamical Astronomy 89(3), (2004) 235–266.

[7] J.F. Cariñena, M.F. Rañada and M. Santander, Central potentials on spaces of constant curvature: The Kepler problem on the two-dimensional sphere \( S^2 \) and the hyperbolic plane \( \mathbb{H}^2 \), J. Math. Phys. 46, (2005), 052702.
[8] F. Diacu, The planar isosceles problem for Manev’s gravitational law, *J. Math. Phys.* 34(12), (1993) 5671–5690.

[9] F. Diacu, Regularization of partial collisions in the n-body problem, *Differential Integral Equations* 5(1), (1992) 103–136.

[10] F. Diacu, Near-collision dynamics for particle systems with quasihomogeneous potentials, *J. Differential Equations* 128, (1996) 58–77.

[11] F. Diacu, On the singularities of the curved n-body problem, *Trans. Amer. Math. Soc.* 363 (2011) 2249–2264.

[12] F. Diacu, Polygonal homographic orbits of the curved n-body problem, *Trans. Amer. Math. Soc.* 364, 5 (2012) 2783–2802.

[13] F. Diacu, Relative equilibria in the 3-dimensional curved n-body problem, *Memoirs Amer. Math. Soc.* 228, 1071 (2013).

[14] F. Diacu, *Relative Equilibria of the Curved N-Body Problem*, Atlantis Studies in Dynamical Systems, vol. 1, Atlantis Press, Amsterdam, 2012.

[15] F. Diacu, The non-existence of centre-of-mass and linear-momentum integrals in the curved n-body problem, [arXiv:1202.4739](http://arxiv.org/abs/1202.4739), 12 p.

[16] F. Diacu, The curved N-body problem: risks and rewards, *Math. Intelligencer* 35, 3 (2013) 24–33.

[17] J. Delgado, F.N. Diacu, E.A. Lacomba, A. Mingarelli, V. Mioc, E. Perez, C. Stoica, The global flow of the Manev problem, *J. Math. Phys.* 37(6), (1996) 2748–2761.

[18] F. Diacu, S. Kordlou, Rotopulsators of the curved N-body problem, *J. Differ. Equations* 255, (2013) 2709–2750.

[19] F. Diacu, V. Mioc, and C. Stoica, Phase-space structure and regularization of Manev-type problems, Nonlinear Analysis 41, 1029–1055 (2000).

[20] F. Diacu and E. Pérez-Chavela, Homographic solutions of the curved 3-body problem, *J. Differ. Equations* 250, (2011) 340–366.

[21] F. Diacu, E. Pérez-Chavela and M. Santoprete, Saari’s conjecture in the collinear case, *Trans. Amer. Math. Soc.* 357, 10 (2005) 4215–4223.
[22] F. Diacu, E. Pérez-Chavela and M. Santoprete, The n-body problem in spaces of constant curvature, arXiv:0807.1747, 54 p.

[23] F. Diacu, E. Pérez-Chavela and M. Santoprete, The n-body problem in spaces of constant curvature. Part I: Relative equilibria, J. Nonlinear Sci. 22, 2 (2012) 247–266.

[24] F. Diacu, E. Pérez-Chavela and M. Santoprete, The n-body problem in spaces of constant curvature. Part II: Singularities, J. Nonlinear Sci. 22, 2 (2012) 267–275.

[25] F. Diacu, E. Pérez-Chavela, M. Santoprete, Central configurations and total collisions for quasihomogeneous n-body problems, Nonlinear Analysis 65, (2006) 1425–1439.

[26] F. Diacu, S. Popa, All the Lagrangian relative equilibria of the curved 3-body problem have equal masses, J. Math. Phys. (to appear).

[27] F. Diacu, B. Thorn, Rectangular orbits of the curved 4-body problem, Proc. Amer. Math. Soc. (to appear).

[28] M. Hampton, R. Moeckel, Finiteness of Relative Equilibria of the four-body problem, Inventiones mathematicae, 163 no. 2 289–312.

[29] W. Killing, Die Rechnung in den nichteuklidischen Raumformen, J. Reine Angew. Math. 89, (1880) 265–287.

[30] W. Killing, Die Mechanik in den nichteuklidischen Raumformen, J. Reine Angew. Math. 98, (1885) 1–48.

[31] W. Killing, Die Nicht-Euklidischen Raumformen in Analytischer Behandlung, Teubner, Leipzig, 1885.

[32] V.V. Kozlov and A.O. Harin, Kepler’s problem in constant curvature spaces, Celestial Mech. Dynam. Astronom. 54, (1992) 393–399.

[33] R.P. Kuz’mina, On an upper bound for the number of central configurations in the planar n-body problem, Sov. Math. Dokl. 18, (1977) 818–821.
[34] H. Liebmann, Die Kegelschnitte und die Planetenbewegung im nichteuklidischen Raum, *Berichte Königl. Sächsischen Gesell. Wiss., Math. Phys. Klasse* 54, (1902) 393–423.

[35] H. Liebmann, Über die Zentralbewegung in der nichteuklidische Geometrie, *Berichte Königl. Sächsischen Gesell. Wisse., Math. Phys. Klasse* 55, (1903) 146–153.

[36] H. Liebmann, *Nichteuklidische Geometrie*, G.J. Göschen, Leipzig, 1905; 2nd ed. 1912; 3rd ed. Walter de Gruyter, Berlin Leipzig, 1923.

[37] J. Llibre, On the number of central configurations in the $N$-body problem, *Celest. Mech. Dyn. Astron.* 50, (1991) 89–96.

[38] N.I. Lobachevsky, The new foundations of geometry with full theory of parallels [in Russian], 1835-1838, In Collected Works, V. 2, GITTL, Moscow, 1949, p. 159.

[39] F.R. Moulton, The Straight Line Solutions of the Problem of $n$ Bodies, *Ann. Math.* 12, (1910) 1–17.

[40] J. Palmore, Collinear relative equilibria of the planar $N$-body problem, *Celest. Mech. Dynam. Astron.* 28, no. 1, (1982) 17–24.

[41] E. Pérez-Chavela, D. Saari, A. Susin and Z. Yan, Central configurations in the charged three body problem, *Contemp. Math.* 198, 137-155 (1996).

[42] E. Pérez-Chavela and L. Vela-Arevalo, Triple collision in the quasihomogeneous collinear three-body problem, *J. Differ. Equations* 148, 186–211 (1998).

[43] G. Roberts, A continuum of relative equilibria in the five-body problem, *Phys. D.* 127, (1999) 141–145.

[44] D. Saari, On the role and properties of n-body central configurations, *Celestial Mech.* 21, (1980) 9-20.

[45] E. Schering, Die Schwerkraft im Gaussischen Räume, *Nachr. Königl. Gesell. Wiss. Göttingen*, 13 July, 15 (1873), 311–321.
[46] E. Schering, Die Schwerkraft in mehrfach ausgedehnten Gaussischen und Riemannschen Räumen, Nachr. Königl. Gesell. Wiss. Göttingen, 26 February, 6 (1873), 149–159.

[47] M. Shub, Relative Equilibria and Diagonals, appendix to S. Smale’s paper ”Problems on the Nature of Relative Equilibria in Celestial Mechanics”, Manifolds, Proc. MUFFIC Summer School on Manifolds, Amsterdam, Springer Lecture Notes in Mathematics 197, (1971), 199–201.

[48] S. Smale, Mathematical problems for the next century, Mathematical Intelligencer 20, (1998) 7–15.

[49] S. Smale, Problems on the nature of relative equilibria in celestial mechanics, Lecture Notes in Mathematics 197, (1971), 194–198.

[50] S. Smale, Topology and Mechanics, II, The planar $n$-body problem, Invent. Math. 11, (1970), 45–64.

[51] P. Tibboel, Polygonal homographic orbits in spaces of constant curvature, Proc. Amer. Math. Soc. 141, (2013), 1465–1471.

[52] P. Tibboel, Existence of a class of rotopulsators, J. Math. Anal. Appl. 404, (2013) 185–191.

[53] P. Tibboel, Existence of a lower bound for the distance between point masses of relative equilibria in spaces of constant curvature, J. Math. Anal. Appl. 416, (2014) 205–211.

[54] A. Wintner, The Analytical Foundations of Celestial Mechanics, Princeton University Press, 1941.