This paper is dedicated to Professor Philippe Bougerol to thank him for many discussions which have been a great source of joy and inspiration to me, and a valuable aid especially in the work presented here.

**Abstract.** We consider the coadjoint action of a Loop group of a compact group on the dual of the corresponding centrally extended Loop algebra and prove that a Brownian motion in a Cartan subalgebra conditioned to remain in an affine Weyl chamber - which can be seen as a space time conditioned Brownian motion - is distributed as the radial part process of a Brownian sheet on the underlying Lie algebra.

1. Introduction

It is a famous result that a real Brownian motion conditioned in Doob’s sense to remain positive, is distributed as a Bessel 3 process, i.e. as the radial process of a 3-dimensional Brownian motion. More generally, if one considers the adjoint action of a compact Lie group on its Lie algebra, the radial part process of a Brownian motion in the Lie algebra is distributed as a Brownian motion in a Cartan subalgebra conditioned to remain in a Weyl chamber. In this paper, we consider the coadjoint action of a Loop group of a compact group on the dual of the corresponding centrally extended Loop algebra and prove that a Brownian motion in a Cartan subalgebra conditioned to remain in an affine Weyl chamber - which can be seen as a space time conditioned Brownian motion - is distributed as the radial part process of a Brownian sheet on the underlying Lie algebra.

Let us be more precise. Let $K$ be a connected compact Lie group, $\mathfrak{k}$ its Lie algebra, and $\mathfrak{t}$ a Cartan subalgebra. One considers a Weyl chamber in $\mathfrak{t}$. Then the orbits of $\mathfrak{t}$ under the adjoint action of $K$ are parametrized by the Weyl chamber. Actually for any $x \in \mathfrak{t}$, it exists a unique vector in the Weyl chamber which is in the same orbit as $x$. This vector is called the radial part of $x$. The Lie algebra $\mathfrak{t}$ is equipped with an $\text{Ad}(K)$-invariant scalar product, whose restriction to the Cartan subalgebra is invariant for the action of the Weyl group. If we consider the radial part process of a standard Brownian motion on $\mathfrak{k}$, it is a classical fact that this process is distributed as its projection on the Cartan subalgebra $\mathfrak{t}$ - which is a Brownian motion on $\mathfrak{t}$ - conditioned in Doob’s sense to remain forever in the Weyl
chamber. The Kirillov’s character formula is at the heart of the connection between the two processes.

An affine Lie algebra can be realized as a central extension of a loop algebra $L\mathfrak{t}$. Considering the coadjoint action of the loop group $L\mathcal{K}$ on the dual of the centrally extended loop algebra, one defines the radial part of an element of this dual algebra (see Pressley and Segal [17]). Frenkel established in [9] a Kirillov character type formula in the framework of affine Lie algebras considering a Gaussian measure on the dual of $L\mathfrak{t}$ - basically a Brownian motion on $\mathfrak{t}$. In his approach, the conditional law of a Brownian motion on $\mathfrak{t}$ given its radial part provides a natural measure on the corresponding orbit under the action of $L\mathcal{K}$. It highly suggested that one could construct a process on a loop algebra, whose the radial part process would be distributed as a projection on a Cartan subalgebra, conditioned to remain in an affine Weyl chamber.

The affine Weyl chamber is a fundamental domain for the action of the Weyl group on the Tits cone. The first difficulty is that there is no Euclidean structure on the Cartan subalgebra of an affine Lie algebra, which would be invariant for the action the Weyl group. So there is no natural Brownian motion to consider on it. The Kirillov’s orbit method, on which we based our intuition, suggests a connection between coadjoint orbits for the action of a Loop group on the dual of an affine Lie algebra and irreducible representations of the affine Lie algebra. Tensor product of irreducible representations of an affine Lie algebra makes appear a drift in the direction of the fundamental weight $\Lambda_0$. The idea is to consider a process with such a drift in the direction of $\Lambda_0$ - which can be seen as a time component - living in an affine Lie algebra. In this paper, we construct such a process, considering a Brownian sheet on the Lie algebra $\mathfrak{k}$ and prove that the corresponding radial part process is distributed as a projection on a Cartan subalgebra, conditioned to remain forever in an affine Weyl chamber.

The paper is organized as follows. In section 2 we describe the orbits for the coadjoint action of a Loop group of a compact group on the dual of the associated centrally extended loop algebra. In particular we define a notion of radial part for this action, which is suitable for our context. In section 4 we briefly recall the necessary background on representation theory of affine Lie algebras. Sections 3 and 5 are basically a reformulation of the main results of [9]. In section 3 we compute the conditional law of a Brownian motion indexed by $[0, 1]$ given the end point of its stochastic exponential. In section 5 we recall how Frenkel proves that this conditional law leads to a Kirillov character formula for affine Lie algebras. In section 6 we introduce a Brownian motion on a Cartan subalgebra of an affine Lie algebra conditioned - in Doob’s sense - to remain forever in an affine Weyl chamber. We prove in section 7 that this conditioned Doob process has the same law as the radial part process of a Brownian sheet on $\mathfrak{t}$.

2. Action of Loop group and its orbits

Loop group and its action. The following presentation is largely inspired by the one given in [17]. Let $K$ be a connected, simply connected, compact Lie group and $\mathfrak{k}$ its Lie algebra. By compactness, without loss of generality, we suppose that $K$ is a matrix Lie group. The adjoint action of $K$ on itself, which is denoted by $\text{Ad}$, is defined by $\text{Ad}(k)(u) = kuk^*$, $k, u \in K$. The induced adjoint action of $K$ on its Lie algebra $\mathfrak{k}$ is still denoted by $\text{Ad}$ and is defined by $\text{Ad}(k)(x) = kxk^*$, $k \in K$, $x \in \mathfrak{k}$. 
The Lie bracket on \( \mathfrak{k} \) is denoted by \([.,.]_{\mathfrak{k}}\). We equipp \( \mathfrak{k} \) with an \( \text{Ad}(K) \)–invariant inner product \( (.,.) \), i.e. the negative of the Killing form. We denote by \( e \) the identity of \( K \). We consider the group \( \text{Map}([0, 1], K) \) of Borel measurable maps from \([0, 1]\) to \( K \), the group law being pointwise composition, and the subgroup \( LK \) of smooth loops from \([0, 1]\) to \( K \),

\[
LK = \{ f : [0, 1] \to K, \text{ } f \text{ is smooth }, f(0) = f(1) \}.
\]

The Lie algebra \( L\mathfrak{k} \) over \( \mathbb{R} \) of \( LK \) is the loop algebra of smooth loops from to \([0, 1]\) to \( \mathfrak{k} \). We define a centrally extended Lie algebra \( L\mathfrak{k} \oplus \mathbb{R}c \) equipped with a Lie bracket given by

\[
[\xi + \lambda c, \eta + \mu c] = [\xi, \eta]_{\mathfrak{k}} + \omega(\xi, \eta)c,\]

for \( \xi, \eta \in L\mathfrak{k}, \lambda, \mu \in \mathbb{R} \), where \( \omega \) is the cocycle defined by \( \omega(\xi, \eta) = \int_0^1 (\xi'(t), \eta(t)) dt \), and \([\xi, \eta]_{\mathfrak{k}} \) is defined pointwise. A Cartan subalgebra of the extended Loop algebra is \( \mathfrak{t} \oplus \mathbb{R}c \), where \( \mathfrak{t} \) is identified with the set of \( \mathfrak{t} \)-valued constant loops. The Lie bracket actually defines a Lie algebra action of \( L\mathfrak{k} \) on \( L\mathfrak{k} \oplus \mathbb{R}c \) given by

\[
\xi.(\eta + \mu c) = [\xi, \eta] + \omega(\xi, \eta)c,
\]

for any \( \xi \in L\mathfrak{k}, (\eta, \mu) \in L\mathfrak{k} \times \mathbb{R} \). This action comes from the adjoint action of \( LK \) on \( L\mathfrak{k} \oplus \mathbb{R}c \) defined by

\[
(1) \quad \gamma.(\lambda + \lambda \Lambda_0) = [\text{Ad}^*(\gamma)\phi - \lambda \int_0^1 (\gamma_s' \gamma_s^{-1}, \cdot) ds] + \lambda \Lambda_0,
\]

for any \( \phi \in (L\mathfrak{k})^*, \lambda \in \mathbb{R} \), where \( \int_0^1 (\gamma_s' \gamma_s^{-1}, \cdot) ds \) stands for the linear form defined by

\[
x \in \mathfrak{k} \mapsto \int_0^1 (\gamma_s' \gamma_s^{-1}, x_s) ds.
\]

Notice that the coadjoint action of the loop group doesn’t affect the level, i.e. the coordinate in \( \Lambda_0 \). Let us equipp \( L\mathfrak{k} \) with the \( L_2 \)-norm, and consider its completion \( L_2([0, 1], \mathfrak{k}) \) with respect to the \( L_2 \)-norm. We define

\[
LH^1([0, 1], K) = \{ \gamma : [0, 1] \to K, \gamma \text{ is absolutely continuous, } \gamma^{-1} \gamma' \in L_2([0, 1], \mathfrak{k}), \gamma(0) = \gamma(1) \}
\]

\[
\mathcal{H}^1([0, 1], \mathfrak{k}) = \{ f \in L_2([0, 1], \mathfrak{k}), f(0) = 0 \},
\]

\[
H^1([0, 1], K) = \{ h : [0, 1] \to K, h \text{ is absolutely continuous, } h^{-1} h' \in L_2([0, 1], \mathfrak{k}), h(0) = e \}.
\]

If we denote by \((L_2([0, 1], \mathfrak{k}))' \) the topological dual of \( L_2([0, 1], \mathfrak{k}) \) then \((2)\) defines an action of \( LH^1([0, 1], K) \) on \((L_2([0, 1], \mathfrak{k}))' \oplus \mathbb{R}\Lambda_0 \), given by

\[
\gamma.(\phi_x + \lambda \Lambda_0) = \int_0^1 (\gamma_x' x_s^{-1}, \cdot) ds - \lambda \int_0^1 (\gamma_s' x_s^{-1}, \cdot) ds + \lambda \Lambda_0,
\]

for any \( \gamma \in LH^1([0, 1], K) \), where \( \phi_x \) is a linear form defined on \( L_2([0, 1], \mathfrak{k}) \) by

\[
\phi_x(y) = \int_0^1 (y_t, x_t') dt,
\]
for any $\gamma \in LH^1([0,1], K)$, $x \in H^1([0,1], \mathfrak{t})$. This action gives rise to an action of $LH^1([0,1], K)$ on $H^1([0,1], \mathfrak{t}) \oplus \mathbb{R}_{\lambda_0}$ defined by

$$\gamma.(x + \lambda \lambda_0) = \int_0^1 (\gamma_x x_s \gamma_s^{-1} - \lambda \gamma_s \gamma_s^{-1}) \, ds + \lambda \lambda_0,$$

for any $\gamma \in LH^1([0,1], K)$, $x \in H^1([0,1], \mathfrak{t})$ and $\lambda \in \mathbb{R}$, which satisfies

$$\phi(\gamma.(x + \lambda \lambda_0) - \lambda \lambda_0) + \lambda \lambda_0 = \gamma.(\phi_x + \lambda \lambda_0).$$

There is another way to make this action appear naturally. One defines a group action of the loop group $LH^1([0,1], K)$, $x \in H^1([0,1], \mathfrak{t})$ and $\lambda \in \mathbb{R}$, which satisfies

$$\phi(\gamma.(x + \lambda \lambda_0) - \lambda \lambda_0) + \lambda \lambda_0 = \gamma.(\phi_x + \lambda \lambda_0).$$

Roots and weights. We choose a set $\Sigma$ of simple roots of $\mathfrak{r}$, $\{\alpha \}$ for any path $x \in H^1([0,1], \mathfrak{t})$ it exists a unique $X \in H^1([0,1], K)$ such that $\lambda dX = X \, dx$. For any $x \in H^1([0,1], \mathfrak{t})$, and $\lambda \in \mathbb{R}$, one defines $\epsilon(x + \lambda \lambda_0)$ as the unique map in $H^1([0,1], K)$ which satisfies this differential equation. As

$$\lambda d(\gamma_0 \epsilon(x + \lambda \lambda_0)) \gamma^{-1} = \gamma_0 \epsilon(x + \lambda \lambda_0) \gamma^{-1} (\gamma \, dx \gamma^{-1} - \lambda \, d\gamma \gamma^{-1}),$$

the action of $LH^1([0,1], K)$ on $H^1([0,1], K)$, and the action of $LH^1([0,1], K)$ on $H^1([0,1], \mathfrak{t}) \oplus \mathbb{R}_{\lambda_0}$, satisfy

$$\epsilon(\gamma.(x + \lambda \lambda_0)) = \gamma.\epsilon(x + \lambda \lambda_0)$$

when $\lambda \neq 0$.

Roots and weights. We choose a maximal torus $T$ of $K$ and denote by $\mathfrak{t}$ its Lie algebra. We denote by $\mathfrak{g}$ the complexification of $\mathfrak{t}$, i.e. $\mathfrak{g} = \mathfrak{t} \oplus i\mathfrak{t}$. We consider the set of real roots

$$R = \{\alpha \in \mathfrak{t}^* : \exists X \in \mathfrak{g} \setminus \{0\}, \forall H \in \mathfrak{t}, [H, X] = 2i\pi \alpha(H) X\}.$$

We choose a set $\Sigma$ of simple roots of $R$ and denote by $R_+$ the set of positive roots. The half sum of positive roots is denoted by $\rho$. Letting for $\alpha \in R$,

$$\mathfrak{g}_\alpha = \{X \in \mathfrak{g} : \forall H \in \mathfrak{t}, [H, X] = 2\pi \alpha(H) X\},$$

the coroot $\alpha^\vee$ of $\alpha \in \Sigma$ is defined to be the only vector of $\mathfrak{t}$ in $[\mathfrak{g}_\alpha, \mathfrak{g}_{-\alpha}]$ such that $\alpha(\alpha^\vee) = 2$. The dual Coxeter number denoted by $h^\vee$ is equal to $1 + \rho(\theta^\vee)$, where $\theta^\vee$ is the highest coroot. We denote respectively by $Q = \sum_i \mathbb{Z} \alpha_i$ and $Q^\vee = \sum_i \mathbb{Z} \alpha_i^\vee$ the root and the coroot lattice. The weight lattice $\{\lambda \in \mathfrak{t}^* : \lambda(\alpha^\vee) \in \mathbb{Z}, \forall \alpha \in \Sigma\}$ is denoted by $P$ and the set $\{\lambda \in \mathfrak{t}^* : \lambda(\alpha^\vee) \in \mathbb{N}, \forall \alpha \in \Sigma\}$ of dominant weights is denoted by $P_+$.

Orbits and radial part of a path from $H^1([0,1], \mathfrak{t})$.

**Proposition 2.1.** Let $x, y \in H^1([0,1], \mathfrak{t})$ and $\lambda \in \mathbb{R}_+^\times$.

1. For any $\gamma \in LH^1([0,1], K)$, one has $\gamma.(x + \lambda \lambda_0) = (y + \lambda \lambda_0)$ if and only if $\gamma.(\phi_x + \lambda \lambda_0) = \phi_y + \lambda \lambda_0$.

2. It exists $\gamma \in LH^1([0,1], K)$, such that $\gamma.(x + \lambda \lambda_0) = y$ if and only if it exists $u \in K$ such that $Ad(u) \epsilon(x + \lambda \lambda_0)_1 = \epsilon(y + \lambda \lambda_0)_1$. 

For the first point comes from identity (3). For the second we write that if $\text{Ad}(u)\epsilon(x + \lambda \Lambda_0)_1 = \epsilon(y + \lambda \Lambda_0)_1$, and $\gamma = \epsilon(y + \lambda \Lambda_0)u\epsilon(x + \lambda \Lambda_0)^{-1}$, then $\gamma \in LH^1([0, 1], K)$ and $\gamma.(x + \lambda \Lambda_0) = y + \lambda \Lambda_0$. $\square$

For $\alpha \in \Sigma$ the fundamental reflection $s_{\alpha^\vee}$ is defined on $t$ by

$$s_{\alpha^\vee}(x) = x - \alpha(x)\alpha^\vee, \quad \text{for } x \in t.$$ 

We consider the extended affine Weyl group generated by the reflections $s_{\alpha^\vee}$ and the translations by $\alpha^\vee$, $x \in t \mapsto x + \alpha^\vee$, for $\alpha \in \Sigma$. The fundamental domain for its action on $t$ is

$$A = \{x \in t : \forall \alpha \in R_+, 0 \leq \alpha(x) \leq 1\}$$

(see for instance section 4.8 of [11]). For $x \in K$, one defines $O_x$ as the adjoint orbit through $x$, i.e.

$$O_x = \{y \in K : \exists u \in K, y = u x u^\ast\}.$$ 

We consider the exponential map $\exp : t \to K$. As $K$ is simply connected, the set of conjugacy classes $K/\text{Ad}(K)$ is in one-to-one correspondence with the fundamental domain $A$, i.e. for all $u \in K$, it exists a unique $x \in A$ such that $u \in O_{\exp(x)}$ (see [4] for instance). Thus given $\lambda \in \mathbb{R}_+^*$, every paths in $H^1([0, 1], t)$ is conjugated to a straight path

$$s \in [0, 1] \mapsto s y,$$

for some $y \in A$, and one can give the following definition for the radial part of $x \in H^1([0, 1], t)$, given a positive level $\lambda \in \mathbb{R}_+^*$.

**Definition 2.2.** For $(x, \lambda) \in H^1([0, 1], t) \times \mathbb{R}_+^*$, one defines the radial part of $x + \lambda \Lambda_0$ as the unique element $r$ in $A$ such that $\epsilon(x + \lambda \Lambda_0)_1 \in O_{\exp(r)}$. It is denoted by $\text{rad}(x + \lambda \Lambda_0)$.

**Radial part of a continuous semi-martingale.** The aim of this part is to define the radial part of a $t$-valued Brownian path. Such a path is not in $H^1([0, 1], t)$ but one can define a stochastic exponential of a Brownian path, which allows us to define its radial part, by analogy with what we’ve done above. Let $(\Omega, \mathcal{F}, (\mathcal{F}_s)_{s \in [0, 1]}, \mathbb{P})$ be a filtered probability space. The following results have been proved for instance in [12] or [15]. If $(x_s)_{s \in [0, 1]}$ is an $t$-valued continuous semi-martingale and $\lambda \in \mathbb{R}_+^*$, then the stochastic differential equation

$$\lambda \, dX = X \circ dx,$$

where $\circ$ stands for the Stratonovitch integral, has a unique solution starting from 0. Such a solution is a $K$-valued process, that we still denote by $(\epsilon(x + \lambda \Lambda_0)_{s})_{s \in [0, 1]}$. This is the Stratonovitch stochastic exponential of $\frac{\lambda}{2}$. Note that given $\lambda \in \mathbb{R}_+^*$, $x$ can be recovered from $\epsilon(x + \lambda \Lambda_0)$ as $dx = \lambda \epsilon(x + \lambda \Lambda_0)^{-1} \circ dt(x + \lambda \Lambda_0)$.

**Definition 2.3.** For $\lambda \in \mathbb{R}_+^*$, and $x = (x_s)_{s \in [0, 1]}$ a continuous $t$-valued semi-martingale, one defines the radial part of $(x + \lambda \Lambda_0)$ as the unique element $r$ in $A$ such that $\epsilon(x + \lambda \Lambda_0)_1 \in O_{\exp(r)}$. It is denoted by $\text{rad}(x + \lambda \Lambda_0)$. 

\[ (4) \]
3. The Conditional Law of a Brownian Motion in $\mathfrak{t}$ Given the End-Point of Its Stochastic Exponential

We denote by $\Lambda$ the kernel of the restriction $\exp_t$ and by $\Lambda^*$ the set of integral weights $\{\lambda \in \mathfrak{t}^* : \lambda(\Lambda) \in \mathbb{Z}\}$, which is included in $P$ since $\alpha^\vee \in \Lambda$ (see [4] for instance). Thus, we define an application $\theta_{\lambda}$ on $T$, when $\lambda \in \Lambda^*$, by letting $\theta_\lambda(\exp(x)) = e^{2\pi i \lambda(x)}$, for $x \in \mathfrak{t}$. The irreducible representations of $K$ are parametrized by the set $\Lambda^*_+ = \Lambda^* \cap C^\vee$, where $C^\vee = \{\lambda \in \mathfrak{t}^* : \lambda(\alpha^\vee) \geq 0, \alpha \in \Sigma\}$. Here, $K$ is supposed to be simply connected, so that $P = \Lambda^*$ and $P^*_+ = \Lambda^*_+$. We denote by $c_{\lambda}$ the character of the irreducible representation with highest weight $\lambda \in P_+$. We recall that $t$ is equipped with a $W$-invariant scalar product $(\cdot, \cdot)$. We identify $t$ and $t^*$ via $(\cdot, \cdot)$ and still denote by $(\cdot, \cdot)$ the scalar product on $t^*$. We write sometimes $e^x$ instead of $\exp(x)$, for $x \in \mathfrak{t}$.

**Brownian motion on $K$.** Let $(\Omega, (\mathcal{F}_s)_{s \in [0,1]}, \mathbb{P})$ be a probability space, where $(\mathcal{F}_s)_{s \in [0,1]}$ is the natural filtration of a $\mathfrak{t}$-valued standard continuous Brownian motion $(x^\sigma)_s \in [0,1]$ defined on $\Omega$, with variance $\sigma > 0$. In this section we consider the stochastic exponential $e^{x^\sigma + \lambda \Lambda_0}$ only for $\lambda = 1$. As for any $\lambda \in \mathbb{R}^*_+$, $e^{x^\sigma + \lambda \Lambda_0} = e^{x^2 \lambda + \Lambda_0}$, there is no loss of generality. In the sequel we let $e^{x^\sigma} = e^{x^2 + \Lambda_0}$. The stochastic exponential $(e^{x^\sigma})_{s \in [0,1]}$ is a left Levy process on $K$ starting from $e$, with transition probability $(p_s^\sigma)_{s \in [0,1]}$ with respect to the Haar measure on $K$, defined on $K$ by

$$p_s^\sigma(x, y) = p_s^\sigma(e, x^{-1} y) = \sum_{\lambda \in P_+} c_{\lambda}(e) c_{\lambda}(x^{-1} y) e^{-\frac{\sigma^2}{2} (||\lambda + \rho||^2 - ||\rho||^2)}.$$

$s \in [0,1], x, y \in K$. In the sequel we write $p_s^\sigma(x)$ instead of $p_s^\sigma(e, x)$. This process is a Brownian motion on $K$. The following proposition states a Girsanov formula for a Brownian motion on a compact Lie group. It is proved for instance in [4], [10] or [12]. For a $\mathfrak{t}$-valued $L_2$ function $y$, and a $\mathfrak{t}$-valued continuous semi-martingale $(x_s)_{s \in [0,1]}$, $\int_0^t (y_s, dx_s)$ is defined as the stochastic integral of $y$ with respect to $x$, for any $t \in [0,1]$.

**Theorem 3.1.** Let $(x^\sigma_s)_{s \in [0,1]}$ be a Brownian motion on $\mathfrak{t}$, with variance $\sigma \in \mathbb{R}^*_+$, and $h \in H^1([0,1], K)$. If $\mu^\sigma$ is the law of $(e^{x^\sigma})_{s \in [0,1]}$, then

$$\frac{d(R_h)_* \mu^\sigma}{d\mu^\sigma} = e^{\frac{1}{2} \int_0^t (h_z^{-1} h'_s, dx^\sigma_s) - \frac{1}{2} \sigma \int_0^t (h_z^{-1} h'_s, h_z^{-1} h'_s) ds},$$

where $(R_h)_* \mu^\sigma$ is the law of $e^{x^\sigma}h$.

For $z \in K$, we write $\mathbb{P}^z$, for the probability defined on $\mathcal{F}_1$ as the conditional probability $\mathbb{P}(\cdot | e^{x^\sigma}_1 = z)$. One has for any $s \in (0,1),$

$$(5) \quad \mathbb{P}^z_{|\mathcal{F}_s} = \frac{p_{t-s}^\sigma(e^{x^\sigma})_{s+1} z}{p_1(z)} \mathbb{P}_{|\mathcal{F}_t}.$$

Note that under $\mathbb{P}^z$, $(x^\sigma_s)_{s \in [0,1]}$ remains a continuous semi-martingale (see for instance [1], theorem 14), so that the stochastic integral $\int_0^t (y_s, dx^\sigma_s)$ is well defined under $\mathbb{P}^z$, for any $\mathfrak{t}$-valued $L_2$ function $y$ and $t \in (0,1)$. Theorem 3.1 implies the following proposition, which appears in proposition (5.2.12) of [9].

**Proposition 3.2.** Let $(x^\sigma_s)_{s \in [0,1]}$ be a standard Brownian motion on $\mathfrak{t}$ with variance $\sigma > 0$ and $y$ be a $\mathfrak{t}$-valued $L_2$ function. If $h \in H^1([0,1], K)$ satisfies $h^{-1} h' = y$ then
for any \( t \in (0, 1) \),
\[
e^{-\frac{\sigma}{2} \int_0^t (y_s, y_s) \, ds} \mathbb{E}(e^{\frac{\sigma}{2} \int_0^t (x_s^\sigma) \, ds} | \epsilon(x^\sigma)_1 = z) = \frac{p_t^\sigma(z h_t^{-1})}{p_t^\sigma(z)}.
\]

**Proof.** For \( t \in (0, 1) \),
\[
\mathbb{E}(e^{\frac{\sigma}{2} \int_0^t (y_s, y_s) \, ds} | \epsilon(x^\sigma)_1 = z) = \mathbb{E}(e^{\frac{\sigma}{2} \int_0^t (x_s^\sigma) \, ds} \frac{p_t^\sigma(z h_t^{-1})}{p_t^\sigma(z)}).
\]

Let \( h \in H^1(K) \) such that \( h^{-1} h' = y \). Theorem 3.1 implies that
\[
\mathbb{E}(e^{\frac{\sigma}{2} \int_0^t (y_s, y_s) \, ds} | \epsilon(x^\sigma)_1 = z) = \mathbb{E}(p_t^\sigma(z h_t^{-1})) \frac{p_t^\sigma(z h_t^{-1})}{p_t^\sigma(z)},
\]
which gives the proposition. \( \square \)

The following lemma has been proved in [9].

**Lemma 3.3.** For \( k_1, k_2 \in K \),
\[
\int_K p_\sigma^\sigma(k_1, uk_2 u^*) \, du = \sum_{\lambda \in P_+} \chi_\lambda(k_1^{-1}) \chi_\lambda(k_2) e^{-\frac{\sigma \langle 2\pi \rangle^2}{2} (||\rho + \lambda||^2 - ||\rho||^2)}.
\]

**Proof.** For any \( \lambda \in P_+ \), the characters \( \chi_\lambda \) satisfies
\[
\int_K \frac{\chi_\lambda(k_1^{-1} uk_2 u^*)}{\chi_\lambda(e)} \, du = \frac{\chi_\lambda(k_1^{-1}) \chi_\lambda(k_2)}{\chi_\lambda(e)}\frac{\chi_\lambda(e)}{\chi_\lambda(e)},
\]
for every \( k_1, k_2 \in K \). Thus
\[
\int_K p_\sigma^\sigma(k_1, uk_2 u^*) \, du = \sum_{\lambda \in P_+} \chi_\lambda(e) \int_K \frac{\chi_\lambda(k_1^{-1} uk_2 u^*) \, du}{\chi_\lambda(e)} e^{-\frac{\sigma \langle 2\pi \rangle^2}{2} (||\rho + \lambda||^2 - ||\rho||^2)}
\]
\[
= \sum_{\lambda \in P_+} \chi_\lambda(k_1^{-1}) \chi_\lambda(k_2) e^{-\frac{\sigma \langle 2\pi \rangle^2}{2} (||\rho + \lambda||^2 - ||\rho||^2)}.
\]
\( \square \)

Previous lemma and Proposition 3.2 imply the following one.

**Proposition 3.4.** Let \( (x^\sigma_s)_{s \in [0, 1]} \) be a standard Brownian motion on \( \mathfrak{k} \) starting from \( 0 \), with variance \( \sigma > 0 \). For \( y \) a \( \mathfrak{k} \)-valued \( L_2 \) function, \( r \in A \), and \( t \in (0, 1) \),
\[
e^{-\frac{\sigma}{2} \int_0^t (y_s, y_s) \, ds} \mathbb{E}(e^{\frac{\sigma}{2} \int_0^t (x_s^\sigma) \, ds} | \operatorname{rad}(x^\sigma) = r) = \frac{1}{p_t^\sigma(e^r)} \sum_{\lambda \in P_+} \chi_\lambda(h_t^{-1}) \chi_\lambda(e) e^{-\frac{\sigma \langle 2\pi \rangle^2}{2} (||\rho + \lambda||^2 - ||\rho||^2)},
\]
where \( h \in H^1([0, 1], K) \) satisfies \( h^{-1} h' = y \).
Proof:
\[
\mathbb{E}(e^{\frac{1}{2} \int_0^T (y, dz) \, ds} | \text{rad}(x^\sigma) = r) = \mathbb{E}(e^{\frac{1}{2} \int_0^T (y, dz) \, ds} | \text{rad}(x^\sigma) = r) \\
= e^{\frac{1}{2} \int_0^T (y, dz) \, ds} \mathbb{E} \left( \frac{p_t^\sigma(e(x^\sigma))}{p_t^\sigma(e(x^\sigma))_1} \right) \, \text{rad}(x^\sigma) = r \\
= e^{\frac{1}{2} \int_0^T (y, dz) \, ds} \int_K \frac{p_t^\sigma(ue^u \epsilon_h^{-1})}{p_t^\sigma(e^u)} \, du,
\]
which proves the proposition thanks to lemma 3.3.

4. Affine Lie algebras and their representations

In this part, we consider an affine Lie algebra whose Cartan matrix is an extended Cartan matrix of the simple finite dimensional complex Lie algebra \( g \). Such an algebra is a nontwisted affine Lie algebra. For our purpose, we only need to consider a realization of its Cartan subalgebra.

4.1. Affine Lie algebras. If \( \{\alpha_1, \ldots, \alpha_n\} \) and \( \{\alpha_1^\vee, \ldots, \alpha_n^\vee\} \) are respectively the sets of simple real roots and coroots of \( K \) previously considered, we let

\[
\mathfrak{h} = \text{span}_\mathbb{C}\{\alpha_0^\vee = c - \theta^\vee, \alpha_1, \ldots, \alpha_n, d\},
\]

and

\[
\mathfrak{h}^* = \text{span}_\mathbb{C}\{\alpha_0 = \delta - \theta, \alpha_1, \ldots, \alpha_n, \Lambda_0\},
\]

where

\[
\alpha_i(d) = \delta_{i0}, \quad \delta(\alpha_i^\vee) = 0, \quad \Lambda_0(\alpha_i^\vee) = \delta_{i0}, \quad \Lambda_0(d) = 0.
\]
The Killing form on \( \mathfrak{t} \) is from now on normalized such that \( (\theta^\vee, \theta^\vee) = 2 \). We consider its restriction to \( \mathfrak{t} \) and extend it to \( \mathfrak{h} \) by \( \mathbb{C} \)-linearity, and by letting

\[
(Cc + Cd, t) = 0, \quad (c, c) = (d, d) = 0, \quad (c, d) = 1.
\]
The following definitions mainly come from chapters 1 and 6 of [14]. The linear isomorphism

\[
\nu: \mathfrak{h} \rightarrow \mathfrak{h}^*, \quad h \mapsto (h, \cdot)
\]
identifies \( \mathfrak{h} \) and \( \mathfrak{h}^* \). We still denote \((\cdot, \cdot)\) the induced bilinear form on \( \mathfrak{h}^* \). We record that

\[
(\delta, \alpha_i) = 0, \quad i = 0, \ldots, n, \quad (\delta, \delta) = 0, \quad (\delta, \Lambda_0) = 1.
\]
Notice that here \( \nu(\theta^\vee) = \theta \) and \( (\theta^\vee, \theta^\vee) = (\theta, \theta) = 2 \). One defines the Weyl group \( \tilde{W} \), as the subgroup of \( GL(\mathfrak{h}^*) \) generated by fundamental reflections \( \mathfrak{s}_\alpha, \alpha \in \Pi \), defined by

\[
\mathfrak{s}_\alpha(\beta) = \beta - \beta(\alpha^\vee)\alpha, \quad \beta \in \mathfrak{h}^*.
\]
Under the identification of \( \mathfrak{h} \) and \( \mathfrak{h}^* \), the action of the affine Weyl group on \( \mathfrak{h} \) is defined by \( w_x = \nu^{-1}wv, x \in \mathfrak{h}, w \in \tilde{W} \). The form \((\cdot, \cdot)\) is \( \tilde{W} \)-invariant. The affine Weyl group \( \tilde{W} \) is the semi-direct product \( W \rtimes \Gamma \) (proposition 6.5 chapter 6 of [14]) where \( \Gamma \) is the group of transformations \( t_\gamma, \gamma \in \nu(Q^\vee) \), defined by

\[
t_\gamma(\lambda) = \lambda + \lambda(c)\gamma - ((\lambda, \gamma) + \frac{1}{2}(\gamma, \gamma)\lambda(c))\delta, \quad \lambda \in \mathfrak{h}^*.
\]
Thus, if \( w \in W, \gamma \in \nu(Q') \), and \( \lambda \in \mathfrak{h}^* \),

\[
wt_{\gamma}(\lambda) = w(\lambda + \lambda(c)\gamma) - ((\lambda, \gamma) + \frac{1}{2}(\gamma, \gamma)\lambda(c))\delta.
\]

**Remark 4.1.** Notice that if one identifies the quotient space \( \mathbb{C}c \oplus t \oplus d/\mathbb{C}c \) with \( t \), then \( \hat{W} \) is identified with the extended affine Weyl group and the fundamental domain of \( \mathbb{C}c \oplus t \oplus d/\mathbb{C}d \) for the action of \( \hat{W} \) is identified with the fundamental domain of \( t \) for the action of the extended affine Weyl group. Actually one has

\[
(d + y) \in \hat{C} \iff y \in A,
\]

where \( \hat{C} \) is the fundamental chamber defined by

\[
\hat{C} = \{ x \in \mathfrak{h} : \forall \alpha \in \hat{\Sigma}, \alpha(x) \geq 0 \}.
\]

The affine Weyl chamber \( \hat{C} \) for the \( A_1^{(1)} \) type is drawn in figure 1 below: it is the area delimited by yellow and orange half-planes. The alcove \( A \) is highlighted in green.

**Figure 1.** The affine Weyl Chamber corresponding to \( A_1^{(1)} \).

### 4.2. Weights, highest-weight modules, characters.

The following definitions and properties mainly come from chapters 9 and 10 of [14]. We denote \( \hat{P} \) (resp. \( \hat{P}_+ \)) the set of integral (resp. dominant) weights defined by

\[
\hat{P} = \{ \lambda \in \mathfrak{h}^* : \langle \lambda, \alpha_i^\vee \rangle \in \mathbb{Z}, i = 0, \ldots, n \},
\]

(resp. \( \hat{P}_+ = \{ \lambda \in \hat{P} : \langle \lambda, \alpha_i^\vee \rangle \geq 0, i = 0, \ldots, n \} \),

where \( \langle ., . \rangle \) is the pairing between \( \mathfrak{h} \) and its dual \( \mathfrak{h}^* \). For \( \lambda \in \hat{P}_+ \) we denote \( V(\lambda) \) the irreducible module with highest weight \( \lambda \). The Weyl-Kac character’s formula (Theorem 10.4, chapter 10 of [14]) states that

\[
\hat{ch}(V(\lambda)) = \frac{\sum_{w \in \hat{W}} \det(w) e^{w(\lambda + \bar{\rho}) - \bar{\rho}}}{\prod_{\alpha \in \hat{R}_+} (1 - e^{-\alpha})^{\text{mult}(\alpha)}},
\]

(7)
where \( \bar{\rho} \in \mathfrak{h}^* \) is chosen such that \( \bar{\rho}(\alpha_i^\vee) = 1 \), for all \( i \in \{0, \ldots, n\} \) and \( \text{mult}(\alpha) \) is the dimension of the root space associated to the root \( \alpha \). In particular

\[
\prod_{\alpha \in R_+} (1 - e^{-\alpha})^{\text{mult}(\alpha)} = \sum_{w \in W} \det(w) e^{w(\bar{\rho})}. 
\]

In the sequel we choose \( \bar{\rho} = h^* \Lambda_0 + \rho \). Writing \( \hat{W} = W \ltimes \Gamma \), the Weyl character’s formula becomes by identity (6)

\[
\hat{\chi}(V(\lambda)) = \sum_{\mu \in \hat{\rho}} \frac{\dim(V(\lambda)_\mu)}{\dim(V(\lambda)_0)} e^{\mu(h)}. 
\]

where the sums run over \( \gamma \in \nu(Q^\vee) \) and \( w \in W \). Letting \( e^{\mu}(h) = e^{\mu(h)} \), \( h \in \mathfrak{h} \), the formal character \( \hat{\chi}(V(\lambda)) \) can be seen as a function defined on its region of convergence. Actually the series

\[
\sum_{\mu \in \hat{\rho}} \dim(V(\lambda)_\mu) e^{\mu(h)} 
\]

converges absolutely for every \( h \in \mathfrak{h} \) such that \( \text{Re}(\delta(h)) > 0 \). We denote by \( \hat{\chi}_\lambda(h) \) its limit (see chapter 11 of [13]). The limiting function \( \hat{\chi}_\lambda \) is analytic on the set

\[
Y = \{ h \in \mathfrak{h} : \text{Re}(\delta(h)) > 0 \}
\]

(see chapter 12 of [14]). For \( y \in \mathfrak{h} \), we let \( \hat{\chi}_y := \hat{\chi}_{\nu(y)} \).

5. A Kirillov character type formula for affine Lie algebras

The Kirillov’s character formula for compact groups. For \( \lambda \in \mathfrak{t}^* \), we denote by \( O_\lambda \) the coadjoint orbit in \( \mathfrak{t}^* \) through \( \lambda \) and by \( \mu_\lambda \) the distribution of \( \text{Ad}(U)\lambda \) where \( U \) is a random variable distributed according the Haar measure on \( K \). When the restriction of \( \lambda \) to \( \mathfrak{t} \) (also denoted by \( \lambda \)) is a dominant weight, the Kirillov’s character formula [13] states that for \( x \in \mathfrak{t} \), one has

\[
\int_{O_\lambda + \rho} e^{2i\pi \beta(x)} \mu_{\lambda + \rho}(d\beta) = j(x) \frac{\hat{\chi}_\lambda(e^x)}{\hat{\chi}_\lambda(0)},
\]

where

\[
j(x) = \frac{\pi(x)}{h(x)}, \quad \pi(x) = \prod_{\alpha \in R_+} (e^{i\pi\alpha(x)} - e^{-i\pi\alpha(x)}), \quad h(x) = \prod_{\alpha \in R_+} 2i\pi\alpha(x).
\]

Definition 5.1. For \( x \in \mathfrak{t} \oplus i\mathfrak{t} \) and \( \lambda \in \mathfrak{t}^* \), one defines

\[
\varphi_x(\lambda) = \frac{1}{h(x)} \sum_{w \in W} (-1)^w e^{w(\lambda, x)}, \quad \varphi_0(\lambda) = \frac{h(\lambda)}{h(\rho)}.
\]

The Weyl’s character formula for compact Lie groups states that

\[
\hat{\chi}_\lambda(e^x) = \frac{\sum_{w \in W} e^{2i\pi(w(\lambda + \rho), x)}}{\sum_{w \in W} e^{2i\pi(w(\rho), x)}} = \frac{1}{\pi(x)} \sum_{w \in W} e^{2i\pi(w(\lambda + \rho), x)},
\]

where
for any \( x \in \mathfrak{t} \). Thus the Kirillov’s character formula becomes, by approximation and analytical continuation,

\[
(8) \quad \int_{\mathcal{O}_\lambda} e^{\beta(x)} \mu_\lambda(d\beta) = \frac{\varphi_x(\lambda)}{\varphi_0(\lambda)},
\]

for any \( x \in \mathfrak{t} \oplus i\mathfrak{t} \) and any \( \lambda \in \mathfrak{t}^* \). This formula gives the Fourier transform of the pushforward measure of the measure \( \mu_\lambda \) under the application \( \phi \in \mathfrak{t}^* \mapsto \phi|_t \in \mathfrak{t}^* \), which is known as the Duistermaat-Heckman measure associated to \( \lambda \).

**A Kirillov-Frenkel character formula for affine Lie algebras.** Let \((x^\sigma)_s\) be a Brownian motion on \( \mathfrak{t} \), with variance \( \sigma > 0 \), and its stochastic exponential \((e(x^\sigma))_s\).

A Kirillov character type formula for affine Lie algebras has been proved by Frenkel in [9]. It can be formulated as in the following theorem.

**Theorem 5.2.** For \( y \in \mathfrak{t} \), and \( \lambda \in \widehat{P}_+ \) such that \( \lambda(c) = k \), one has

\[
\mathbb{E}(e^{\frac{1}{2}(x^\sigma,y)}|e(x^\sigma)_1 \in \mathcal{O}_{\exp(\sigma z)}) = \tilde{j}(y) \frac{\chi_\Lambda(d + y)}{\chi_\Lambda(d)}.
\]

where \( \nu(z) = \rho + \lambda - k\Lambda_0 \), \( \sigma = \frac{1}{e^{2\pi i}} \), and

\[
\tilde{j}(y) = \prod_{\alpha \in R^+_+ \setminus \Pi} \left( 1 - e^{-\alpha(d + y)} \right)^{\text{mult}(\alpha)} \prod_{\alpha \in R^+_+} \frac{2i\pi(1 - e^{-\alpha(y)})}{1 - e^{-2\pi\alpha(y)}}.
\]

**Definition 5.3.** For \( x \in \mathfrak{t}, \; y \in \mathfrak{t} \oplus i\mathfrak{t}, \; a, b \in \mathbb{R}^*_+ \), we let

\[
(9) \quad \widehat{\varphi}_{bd + y}(a, x) = \frac{1}{\pi(-\frac{i}{2}y)} \sum_{w \in \hat{W}} (-1)^w e^{(w(ad + x), db + y)}
\]

As \( \hat{W} \) is the semi-direct product \( W \ltimes \Gamma \) one has also

\[
(10) \quad \widehat{\varphi}_{bd + y}(a, x) = \frac{1}{\pi(-\frac{i}{2}y)} \sum_{\gamma \in Q^\vee} \sum_{w \in \hat{W}} (-1)^w e^{(w(x + a\gamma), y)} e^{-b((x, \gamma) + \frac{1}{2}(\gamma, \gamma)a)},
\]

**Remark 5.4.** The previous definition seems to be not valid for \( y \) such that \( \pi(\frac{1}{2}y) = 0 \). Actually, in proposition 5.6 we give another expression for \( \widehat{\varphi}_{bd + y} \), which shows that it is actually well defined for every \( y \in \mathfrak{t} \oplus i\mathfrak{t} \).

The Poisson formula implies the following lemma.

**Lemma 5.5.** For \( x \in \mathfrak{t} \), and \( t \in \mathbb{R}^*_+ \), one has

\[
\sum_{\mu \in P} e^{2i\pi \mu(x) - \frac{1}{2}(2\pi)^2(\mu, \mu)} = C \left( \frac{1}{2\pi t} \right)^{n/2} \sum_{z \in Q^\vee} e^{-\frac{1}{4}(x + z, x + z)}.
\]

where \( C \) is a positive constant independent of \( x \) and \( t \).

**Proof.** As \( \mu \) is a weight in \( P \) if and only if \( \mu(Q^\vee) < Z \), the Poisson summation formula gives for every bounded continuous function \( f \) defined on \( \mathfrak{t}^* \),

\[
\sum_{\mu \in P} f(\mu) = C \sum_{z \in Q^\vee} \hat{f}(z),
\]

where \( \hat{f}(z) = \int_{\mathfrak{t}^*} e^{2i\pi \mu(z)} f(\mu) d\mu \), and \( C \) is a positive constant independent of \( f \). The lemma follows taking \( f \) defined by \( f(\mu) = e^{2i\pi \mu(x) - \frac{1}{2}(2\pi)^2(\mu, \mu)} \), \( \mu \in \mathfrak{t}^* \).
Proposition 5.6. For $x, y \in \mathfrak{t}$, and $\sigma \in \mathbb{R}_+^*$, one has
\[
\hat{\phi}_{d+y}(\frac{1}{\sigma}, \frac{x}{\sigma}) = C \left( \frac{1}{2\pi\sigma} \right)^{-n/2} e^{\frac{1}{2\pi} (y, y) + \frac{1}{2\sigma} (x, x)} \sum_{\mu \in P_+} \pi(x) \chi_\mu(e^x) \chi_\mu(e^{-y}) e^{-\frac{z}{2}(2\pi)^2 ||\mu + \rho||^2},
\]
where $C$ is a positive constant independent of $x$, $y$ and $\sigma$.

Proof. One has to prove that
\[
\frac{1}{\pi(-y)} \sum_{w \in W} (-1)^w e^{\frac{z}{2}(w(d+y),x)}
\]
is equal to
\[
C \left( \frac{1}{2\pi\sigma} \right)^{-n/2} e^{\frac{1}{2\pi} (y, y) + \frac{1}{2\sigma} (x, x)} \sum_{\mu \in P_+} \pi(x) \chi_\mu(e^x) \chi_\mu(e^{-y}) e^{-\frac{z}{2}(2\pi)^2 ||\mu + \rho||^2},
\]
where $C$ is a positive constant independent of $x$, $y$ and $\sigma$. On the one hand, on has
\[
C \sum_{\mu \in P, w \in W} (-1)^w e^{2\pi \mu(x-wy)} e^{-\frac{1}{2}(2\pi)^2 \sigma(x, x)} = \left( \frac{1}{2\pi\sigma} \right)^{n/2} \sum_{z \in \mathbb{Q}^\vee, w \in W} (-1)^w e^{\frac{z}{2}(w(y, x) - \frac{1}{2}(x, x))}
\]
\[
= \left( \frac{1}{2\pi\sigma} \right)^{n/2} e^{-\frac{z}{2}(y, y) - \frac{z}{2}(x, x)} \sum_{w \in W, z \in \mathbb{Q}^\vee} (-1)^w e^{\frac{z}{2}(w(y, x) - \frac{1}{2}(x, x))}.
\]
On the other hand,
\[
\sum_{\mu \in P, w \in W} (-1)^w e^{2\pi \mu(x-wy)} = \sum_{\mu \in P_+, w, \tilde{w} \in W} (-1)^w e^{2\pi \chi_\mu(\frac{1}{\sigma} (-\frac{1}{\sigma}(w(y, x) - \frac{1}{2}(x, x)))}
\]
\[
= \sum_{\mu \in P_+} \chi_\mu(e^x) \pi(x) \chi_\mu(e^{-y}) \pi(-y) e^{-\frac{1}{2}(2\pi)^2 ||\mu + \rho||^2}
\]
\[ \square \]

Remark 5.7. By analytical continuation, the identity remains true for $y \in \mathfrak{t} \oplus i \mathfrak{t}$, $\sigma \in \mathbb{C}$, Re$(\sigma) > 0$.

Propositions 5.6 and 3.4 imply the following theorem, which implies the Kirillov character type formula of theorem 5.2.

Theorem 5.8. For $y \in \mathfrak{L}_2([0, 1], \mathfrak{t})$, and $z \in \mathfrak{A}$, one has for $t \in (0, 1)$,
\[
\mathbb{E}(e^{\frac{1}{\hbar} \int_0^t(y_s, dx_s^e)} | \text{rad}(x^e) = z) = e^{\frac{1}{2\hbar} \int_0^t(y_s, y_s) ds} e^{-\frac{1}{2\hbar} (a, a) \hat{\phi}_{d+a}(\frac{1}{\sigma}, \frac{z}{\sigma})}\frac{\hat{\phi}_{d+a}(\frac{1}{\sigma}, \frac{z}{\sigma})}{\hat{\phi}_{d}(\frac{1}{\sigma}, \frac{z}{\sigma})},
\]
where $\hat{h}^{-1} h'$ and $h_t \in \mathfrak{O}_{e^e}$, with $a \in \mathfrak{t}$. In particular for $y \in \mathfrak{t}$, one has
\[
\mathbb{E}(e^{\frac{1}{\hbar} (y, dx_t^e)} | \text{rad}(x^e) = z) = \frac{\hat{\phi}_{d+y}(\frac{1}{\sigma}, \frac{z}{\sigma})}{\hat{\phi}_{d}(\frac{1}{\sigma}, \frac{z}{\sigma})}.
\]

Remark 5.9. Let us make some non rigorous remarks about these formulae. Previously we have considered the identification of $L_2([0, 1], \mathfrak{t})' \oplus \mathfrak{R} \mathfrak{A}_0$ with $H_1([0, 1], \mathfrak{t})' \oplus \mathfrak{R} \mathfrak{A}_0$ letting $\varphi_x = \int_0^1 (x_s, dx_s^e) ds$ for $x \in H_1([0, 1], \mathfrak{t})$. In the first formula the stochastic integral $\int_0^1 (\cdot, dx_s^e)$ can be seen as a random linear functional. Its conditional law given $e(x^e) \in \mathfrak{O}_{e(z)}$ has to be thought as a measure on a coadjoint orbit through
positive true martingale under $W$ Brownian motion, $\tau$ under which $z$ for $y$ is also a true martingale. Continuity in process (where $C$ is a positive constant, which ensures in particular that $\varphi_{x,u}$ is a true martingale under $W$ Brownian motion, $\tau$ being the time component. Let us consider the stopping time

$$T = \inf\{t \geq 0 : x_t \notin \tilde{C} \} = \inf\{t \geq 0 : b_t \notin \tau_t A\}.$$ 

Lemma 6.1. Let $u > 0$ and $(b_t)_{t \geq 0}$ be a standard Brownian motion on $t$. For $y \in t \oplus it$, $(e^{-\frac{y^2}{2}} \hat{\varphi}_{d+y}(\tau_t, b_t), t \geq 0)$ is a true martingale under $\mathbb{W}_{x,u}$.

Proof. Let us suppose that $\pi(y) \neq 0$. Choose an orthonormal basis $v_1, \ldots, v_n$ of $t$ and consider for $w \in \tilde{W}$ a function $g_w$ defined on $\mathbb{R}_+^* \times \mathbb{R}^n$ by

$$g_w(t, x_1, \ldots, x_n) = e^{(td+x, w(d+y))} e^{\frac{i}{2}(y, y)},$$

where $x = x_1v_1 + \cdots + x_nv_n$. Letting $\Delta = \sum_{i=1}^n \partial_{x_i}^2$, the function $g_w$ satisfies

$$(\frac{1}{2} \Delta + \partial_t)g_w = 0,$$

which implies that $(g_w(t + u, b_t))_{t \geq 0}$ is a local martingale. Its quadratic variation is easily shown to be integrable, so that, it is a true martingale. The sum $\sum_{w \in \tilde{W}} (-1)^w g_w(t + u, b_t)$ converges for the $L^1$ norm, which implies that

$$(e^{\frac{i}{2}(y, y)} \varphi_{d+y}(t+u, b_t), t \geq 0),$$

is also a true martingale. Continuity in $y$ ensures that lemma is true for any $y \in t \oplus it$. \hfill $\square$

Lemma 6.2. Let $t > 0$. If $td + x \in \tilde{C}$, i.e. $x \in tA$, then $\varphi_{d}(t, x) \geq 0$, with equality if and only if $td + x$ is on the boundary of $\tilde{C}$, i.e. $x$ is on the boundary of $tA$.

Proof. Proposition 5.6 implies that for any $z \in t$,

$$\varphi_d(\frac{1}{\sigma}, \frac{z}{\sigma}) = C \times p_\sigma^d (e^z) \pi(z) e^{\frac{1}{2}(2\pi)^2 ||z||^2 - \frac{1}{2}(2\pi)^2 ||\sigma||^2},$$

where $C$ is a positive constant, which ensures in particular that

$$\varphi_d(t, z) \geq 0,$$

for $z \in tA$, with equality if and only if $z \in t\partial A$. \hfill $\square$

Let $(\mathcal{F}_t)_{t \geq 0}$ be the natural filtration of $(x_t)_{t \geq 0}$. As $(\varphi_d(\tau_{tA}, b_{tA}), t \geq 0)$, is a positive true martingale under $\mathbb{W}_{x,u}$ such that $\varphi_d(T, b_T) = 0$, one defines a measure $Q_{x,u}$ on $\mathcal{F}_\infty$ as below.
Lemma 7.1. Let $t > 0$ and $x \in \mathfrak{t}$ such that $ud + x$ is in the interior of $\hat{C}$, i.e. for any $y \in \mathfrak{t} \oplus \mathfrak{i}$, one defines a probability $Q_{u,x}$ letting

$Q_{u,x}(B) = \mathbb{E}_{\mathfrak{W}_{u,x}}(\frac{\hat{\varphi}_d(\tau_{t}, b_t)}{\varphi_d(u, x)} 1_{T \geq t}, B)$, for $B \in \mathcal{F}_t$, $t \geq 0$.

Actually this conditioned Doob process can be obtained as a limit of a space time Brownian motion starting in the affine Weyl chamber, with drift within the affine Weyl chamber, conditioned to remain forever in it, when the drift goes to zero (see [7] and [8] for more details). Lemma 6.1 and the fact that for $y \in \mathfrak{t} \oplus \mathfrak{i}$, $\hat{\varphi}_{d+y}(\tau_T, b_T) = 0$ imply immediately the following proposition.

Proposition 6.4. For $r, t, u \in \mathbb{R}_+^*$, $x$ in the interior of $u \mathfrak{a}$, and $y \in \mathfrak{t} \oplus \mathfrak{i}$, one has

$$E_{Q_{u,x}}(\frac{\hat{\varphi}_{d+y}(\tau_{r+t}, b_{t+r})}{\varphi_d(\tau_{r+t}, b_{t+r})}) = \frac{\hat{\varphi}_{d+y}(u, x)}{\varphi_d(u, x)} e^{(u, y)_t} t.$$ (12)

and

$$E_{Q_{u,x}}(\frac{\hat{\varphi}_{d+y}(\tau_{r+t}, b_{t+r})}{\varphi_d(\tau_{r+t}, b_{t+r})}) \mid \mathcal{F}_r = \frac{\hat{\varphi}_{d+y}(\tau_{r}, b_{r})}{\varphi_d(\tau_{r}, b_{r})} e^{(u, y)_t}.$$ (13)

7. Conditioned space time Brownian motion, and radial part of a Brownian sheet

In this last section we prove that the conditioned Doob process in an affine Weyl chamber previously introduced, has the same law as the radial part process of a Brownian sheet on $\mathfrak{t}$. It is stated in theorem [7.7] Let $(x'_s)_{s \in [0,1], t \in \mathbb{R}_+}$ be a standard Brownian sheet on $\mathfrak{t}$, i.e. for any $t, t' \in \mathbb{R}$, $(x'_s + t - x'_s)_{s \in [0,1]}$ is a $\mathfrak{t}$-valued random process independent of $\sigma(x'_s, s \in [0,1], r \leq t)$, having the same law as $(x'_s)_{s \in [0,1]}$, which is a standard Brownian motion on $\mathfrak{t}$, with variance $t'$. In the sequel we choose a continuous version of it. Proposition [7.3] and corollary [7.4] prove the existence of an entrance law for the conditioned process in the affine Weyl chamber introduced in section [4] and the entrance point $0$.

Lemma 7.1. For any $y \in \mathfrak{t} \oplus \mathfrak{i}$,

$$\frac{\hat{\varphi}_{d+y}(u, x)}{\varphi_d(u, x)}$$

converges towards 1 when $(u, x)$ goes to 0 within the affine Weyl chamber.

Proof.\[\hat{\varphi}_{d+y}(u, x) = c(\frac{u}{2\pi})^{-n/2}e^{\frac{x^2}{2}(y, y)+\frac{1}{2}(x, x)\pi} \sum_{\mu \in \mathcal{P}_+} \chi_\mu(e^{2i}) \chi_\mu(e^{-y}) e^{-\frac{t}{2}(2x)^2||u+\rho||^2},\]

Lemma 13.13 of [14] implies that the dominant term in the sum of the right hand side of the identity is $e^{-\frac{t}{2}(2x)^2||\rho||^2}$, and proposition follows. □

Lemma 7.2. Let us fix $t > 0$. If $\mu_{x,u}$ is the law of $\frac{b_t}{t+u}$ under $Q_{x,u}$ then for any $y \in \mathfrak{t} \oplus \mathfrak{i}$,

$$\int_{\mathfrak{a}} \frac{\hat{\varphi}_{d+y}(t, tz)}{\varphi_d(t, tz)} \mu_{x,u}(dz)$$

converges towards $e^{(x, y)_t}$ when $(x, u)$ goes to 0 within the affine Weyl chamber.
Proof. If $\mu_{x,u}$ is the law of $\frac{b_{x,u}}{\tau_{x,u}}$ under $\mathbb{Q}_{x,u}$ then
\[
\int_A \frac{\hat{\varphi}_{d+y}(t+u)(t+u)z}{\hat{\varphi}_d(t+u)(t+u)z} \, d\mu_{x,u}(z) = \frac{\hat{\varphi}_{d+y}(x,u)}{\hat{\varphi}_d(x,u)} e^{\hat{\gamma}(y,u)}.
\]

Proposition 5.6 and the fact that $|\chi_\mu(e^z)| \leq \frac{\lambda(u+\rho)}{\tau(u)}$ imply that $\frac{\hat{\varphi}_{d+y}(t+u)(t+u)z}{\hat{\varphi}_d(t+u)(t+u)z}$ converges to $\frac{\hat{\varphi}_d(t,tz)}{\hat{\varphi}_d(t,tz)}$, uniformly in $z \in A$, when $u$ goes to 0. Thus lemma follows from lemma 7.1.

\[\square\]

**Proposition 7.3.** Let $(x^t_s)_{s \in [0,1]}$ be a standard Brownian sheet on $\mathfrak{t}$ and $t > 0$. Under $\mathbb{Q}_{u,x}$, $\frac{b_{x,u}}{\tau_{x,u}}$ converges in law towards the radial part of $(\frac{1}{t}x^t_s)_{s \in [0,1]}$, when $(x,u)$ goes to 0 within the affine Weyl chamber.

Proof. Let $t > 0$. By analytical continuation, the second identity of theorem 5.8 remains valid for any $y \in \mathfrak{t} \oplus it$. As $(\frac{1}{t}x^t_s)_{s \in [0,1]}$ is a standard Brownian motion on $\mathfrak{t}$ with variance $\frac{1}{t}$, theorem 5.8 gives for $z \in A$, $y \in \mathfrak{t} \oplus it$,

\[
\mathbb{E}(e^{(x^t_s,y)}) \operatorname{rad}((\frac{1}{t}x^t_s)_s) = z = \frac{\hat{\varphi}_{d+y}(t,tz)}{\hat{\varphi}_d(t,tz)}.
\]

In particular the law $\mu$ of $\operatorname{rad}((\frac{1}{t}x^t_s)_s)$ satisfies, for any $y \in \mathfrak{t} \oplus it$,

\[
\int_A \frac{\hat{\varphi}_{d+y}(t,tz)}{\hat{\varphi}_d(t,tz)} \, d\mu(z) = \mathbb{E}(e^{(x^t_s,y)}) = e^{\frac{1}{t}(y,u)}.
\]

Proposition 5.6 and the fact that $|\chi_\mu(e^z)| \leq \frac{\lambda(u+\rho)}{\tau(u)}$ imply that

\[
z \in A \mapsto \frac{\hat{\varphi}_d(t,zt)}{\pi(z)e_t(z)},
\]

where $e_t(z) = \sum_{\gamma \in \nu(Q^+)} e^{-t(\gamma(z) + \frac{1}{2}(\gamma,\gamma))}$ is smooth on $A$. Thus the Peter-Weyl theorem ensures that for any smooth function $u$ defined on $A$, and $z \in A$,

\[
u(z) \frac{\hat{\varphi}_d(t,zt)}{\pi(z)e_t(z)} = \sum_{\lambda \in P_+} c_\lambda \chi_\lambda(z),
\]

where $c_\lambda = \int_A u(z) \frac{\hat{\varphi}_d(t,zt)}{\pi(z)e_t(z)} \chi_\lambda(z) \pi(z)^2(z) \, dz$, and the convergence stands uniformly and absolutely. Actually $\lim_{|\lambda| \to \infty} \langle \lambda,\lambda \rangle^n c_\lambda = 0$ for all $n \in \mathbb{N}$. As $\hat{\varphi}_d(t,zt)/(\pi(z)e_t(z))$ remains positive on $A$,

\[
u(z) = \sum_{\lambda \in P_+} c_\lambda \chi_\lambda(z) \frac{\pi(z)e_t(z)}{\hat{\varphi}_d(t,zt)}.
\]

As $|\chi_\lambda(z)| \leq d(\lambda + \rho)$ and $\pi(.)e(.)/\hat{\varphi}_d(t,t.)$ is bounded on $A$, the uniform and absolute convergence implies that for any probability measure $\nu$ on $A$,

\[
\int_A u(z)\nu(dz) = \sum_{\lambda \in P_+} c_\lambda \int_A \chi_\lambda(z) \frac{\pi(z)e_t(z)}{\hat{\varphi}_d(t,zt)} \nu(dz).
\]

Expression 10 for the definition of $\hat{\varphi}_{d+y}$ gives,

\[
\hat{\varphi}_{d+2\pi \frac{i}{t}(\lambda+\rho)}(t,tz) = \frac{1}{\pi(-2\pi \frac{i}{t}(\lambda+\rho))} \chi_\lambda(z) \pi(z)e_t(z),
\]

Expression 10 for the definition of $\hat{\varphi}_{d+y}$ gives,
lemma 7.2 implies that
\[
\nu = \text{Identity (14) becomes for } \square \text{ follows from propositions 7.3 and 5.6.}
\]

for any \( \lambda \in P_+ \). Thus one has
\[
\int_A \chi_\lambda(z) \frac{\pi(z)e(z)}{\hat{\varphi}_d(t,zt)} d\mu(z) = \pi(-2\pi i t (\lambda + \rho)) \int_A \frac{\hat{\varphi}_{d+2\pi i (\lambda + \rho)}(t,zt)}{\hat{\varphi}_d(t,zt)} d\mu(z)
\]

\[
= \pi(-2\pi i t (\lambda + \rho)) e^{-\frac{t}{2}(2\pi)^2(\lambda + \rho, \lambda + \rho)},
\]

and identity \([14]\) becomes for \( \nu = \mu, \)
\[
\int_A u(z) d\mu(z) = \sum_{\lambda \in P_+} c_\lambda \pi(-2\pi i t (\lambda + \rho)) e^{-\frac{t}{2}(2\pi)^2(\lambda + \rho, \lambda + \rho)}.
\]

Identity \([14]\) becomes for \( \nu = \mu_{x,u} \)
\[
\int_A u(z) d\mu_{x,u}(z) = \sum_{\lambda \in P_+} c_\lambda \int_A \chi_\lambda(z) \frac{\pi(z)e(z)}{\hat{\varphi}_d(t,zt)} d\mu_{x,u}(z).
\]

As
\[
\int_A \chi_\lambda(z) \frac{\pi(z)e(z)}{\hat{\varphi}_d(t,zt)} d\mu_{x,u}(z) = \pi(-2\pi i t (\lambda + \rho)) \int_A \frac{\hat{\varphi}_{d+2\pi i (\lambda + \rho)}(t,zt)}{\hat{\varphi}_d(t,zt)} d\mu(z)
\]

lemma 7.2 implies that
\[
\int_A u(z) d\mu_{x,u}(z) \text{ converges towards } \int_A u(z) d\mu(z),
\]

as \((x, u)\) goes to 0 within the affine Weyl chamber, which proves the proposition. \(\square\)

**Corollary 7.4.** For any \( t > 0, \)
\[
\lim_{(x,u) \to 0} \mathbb{Q}_{x,u}(b_t \in dz) = C_t \hat{\varphi}_d(t,z) \pi(\frac{z}{t}) 1_A(\frac{z}{t}) W_0(b_t \in dz),
\]
when \((x, u)\) goes to 0 within the affine Weyl chamber.

**Proof.** The Weyl integration formula implies that the density of the radial part of \( \frac{1}{t} x^t \) is equal to \( p_1^t(\pi(z)^2 1_A(z), \) up to a multiplicative constant. Thus corollary follows from propositions 7.3 and 5.6 \(\square\)

One defines a law \( \mathbb{Q}_{0+} \) on \( \sigma(x_u : u > 0) \) letting for \( B \in \sigma(x_u : u \geq t), t > 0, \)
\[
\mathbb{Q}_{0+}(B) = \mathbb{P}_0(C_t \hat{\varphi}_d(t,b_t) \pi(\frac{b_t}{t}) Q_{x_t}(\theta_t B)),
\]

where \( \theta \) is the shift operator.

**Remark 7.5.** Under \( \mathbb{Q}_{0+}, \)
\[
\frac{b_t}{t} \text{ is equal in law to } \text{rad}(\frac{x^t}{t}) \text{ for any } t > 0. \]

Thus for any \( t > 0, \)
\[
\frac{b_t}{t} \text{ is equal in law to } \text{rad}(x^t). \]

As \( \epsilon(x^t)_1 \) converges towards the Haar measure on \( K \) when \( t \) goes to 0, one obtains by the Weyl integration formula that
\[
\lim_{t \to 0} \mathbb{Q}_{0+}(\frac{b_t}{t} \in dx) = C \pi(x)^2 1_A(x) dx,
\]

which can be also deduced from \([15]\).

Let us denote by \( C([0, 1], \mathcal{T}) \) the set of continuous maps from \([0, 1] \) to \( \mathcal{T} \).
Proposition 7.6. Let \((x^t_s)_{s \in [0,1], t \in \mathbb{R}^+}\) be a standard Brownian sheet on \(\mathfrak{t}\). Let \(\Lambda\) be a kernel on \((\mathbb{R}^*_+ \times A) \times (\mathbb{R}^*_+ \times C([0,1], \mathfrak{t}))\), such that for any \((t,y) \in \mathbb{R}^*_+ \times A\), \(\Lambda((t,y), (t,\cdot))\) is the law of a brownian motion \((x^t_s)_{s \in [0,1]}\) on \(\mathfrak{t}\), given \(\text{rad} \left( \frac{1}{t} x^t \right) = y\). Let \((P_t)_{t \in \mathbb{R}^*_+}\) be the transition probability of the Markov process \((\tau_t, \frac{b_t}{t})_{t \in \mathbb{R}^*_+}\) under \(Q_{0^+}\), and \((Q_t)_{t \in \mathbb{R}^*_+}\) the transition probability of \((x^t, t)_{t \in \mathbb{R}^*_+}\). Then for any \(t > 0\) and \((u_0, x_0) \in \mathbb{R}^*_+ \times A\),

\[
P_t \Lambda((u_0, x_0), \cdot) = \Lambda Q_t((u_0, x_0), \cdot).
\]

Proof. It is sufficient to prove that for every measurable function \(y \in L_2([0,1], \mathfrak{r})\), \((u_0, x_0) \in \mathbb{R}^*_+ \times A\), and \(r \in (0,1)\),

\[
\int P_t \Lambda((u_0, x_0), (t + u_0, d z)) e^{\int_0^r (y_z, dz)} = \int \Lambda Q_t((u_0, x_0), (t + u_0, d z)) e^{\int_0^r (y_z, dz)}.
\]

Let \(y \in L_2([0,1], \mathfrak{r})\), \((u_0, x_0) \in \mathbb{R}^*_+ \times A\), \(r \in (0,1)\), \(h \in H^1([0,1], \mathfrak{r})\) such that \(h^{-1} h' = y\), and \(a \in \mathfrak{t}\) such that \(h_r \in \mathcal{O}_{\mathfrak{a}_r}\). On the one hand, one has

\[
\int P_t \Lambda((u_0, x_0), (t + u_0, d z)) e^{\int_0^r (y_z, dz)}
= \int P_t((u_0, x_0), (u_0 + t, d z)) \Lambda((u_0 + t, z), (u_0 + t, d z)) e^{\int_0^r (y_z, dz)}
= \int P_t((u_0, x_0), (u_0 + t, d z)) \frac{e^{\int_0^{t+u_0} f^r_0(y_{s}, y_s) ds - \frac{1}{2} (t+u_0)(a,a) \hat{Q}_{d+a}(t+u_0, (u_0 + t)z) \hat{Q}_{d}(t+u_0, (u_0 + t)z)}}{\hat{Q}_{d}(t+u_0, (u_0 + t)z)}
= e^{\frac{1}{2} (t+u_0) \int_0^r (y_s, y_s) ds - \frac{1}{2} (t+u_0)(a,a) \hat{Q}_{d+a}(u_0, u_0) \hat{Q}_{d}(u_0, u_0)}.
\]

On the other hand,

\[
\int \Lambda((x, u_0), (d z, u_0)) Q_t((z, u), (d z, u + t)) e^{\int_0^r (y_z, dz)}
= e^{\frac{1}{2} \int_0^r (y_s, y_s) ds} \int \Lambda((x, u), (d z, u)) e^{\int_0^r (y_z, dz)}
= e^{\frac{1}{2} \int_0^r (y_s, y_s) ds} e^{\frac{1}{2} (t+u_0)(a,a) \hat{Q}_{d+a}(u, u) \hat{Q}_{d}(u, u) / \hat{Q}_{d}(u, u)}.
\]

For the following theorem we consider a continuous version of the doubly indexed process \((e^{\frac{r}{t}} z^t_s)_{s \in [0,1], t \in \mathbb{R}^*_+}\) (see for instance [6] and references therein for the existence of a continuous version).

Theorem 7.7. Let \((x^t_s)_{s \in [0,1], t \in \mathbb{R}^*_+}\) be a standard Brownian sheet. Under \(Q_{0^+}\), the process \((\frac{b_t}{t})_{t \in \mathbb{R}^*_+}\) is equal in law to the process \((\text{rad} \left( \frac{1}{t} x^t \right))_{t \in \mathbb{R}^*_+}\).

Proof. Applying the criterion described by Rogers and Pitman in [6], the theorem is a consequence of the proposition 7.3 and the intertwining relation established in proposition 7.6.
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