Near-Unity Light Absorption in a Monolayer WS$_2$ Van der Waals Heterostructure Cavity

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ABSTRACT: Excitons in monolayer transition-metal-dichalcogenides (TMDs) dominate their optical response and exhibit strong light–matter interactions with lifetime-limited emission. While various approaches have been applied to enhance light-exciton interactions in TMDs, the achieved strength have been far below unity, and a complete picture of its underlying physical mechanisms and fundamental limits has not been provided. Here, we introduce a TMD-based van der Waals heterostructure cavity that provides near-unity excitonic absorption, and emission of excitonic complexes that are observed at ultralow excitation powers. Our results are in full agreement with a quantum theoretical framework introduced to describe the light–exciton–cavity interaction. We find that the subtle interplay between the radiative, nonradiative and dephasing decay rates plays a crucial role, and unveil a universal absorption law for excitons in 2D systems. This enhanced light–exciton interaction provides a platform for studying excitonic phase-transitions and quantum nonlinearities and enables new possibilities for 2D semiconductor-based optoelectronic devices.

KEYWORDS: TMD Excitons, Unity absorption, 2D materials, Light-matter interaction, Exciton complexes

The remarkable properties of excitons in monolayer TMDs, together with the ability to readily control their charge carrier density, have attracted a significant amount of interest in recent years. This has led to the observation of numerous phenomena, such as higher-order exciton complexes, coupled spin-valley physics, single photon quantum emitters, together with monolayer semiconductor-based lasers, light-emitting-diodes, and photodetectors. Excitons in monolayer TMDs exhibit strong interaction with light, both in absorption and photoemission processes, which facilitates large photovoltaic response enabled by strong peaks in the joint density of states and strong-coupling to name two examples. Unlike their counterparts in quantum-well semiconductors, excitons in TMDs practically dominate the optical response of the material. This stems from their large binding energies, which are a result of the strong Coulomb interaction and reduced screening that arise from their low dimensionality.

The existence of robust excitonic states deep within the bandgap results in an optical bandgap that differs significantly from the electronic one and, thus, prevails over the standard electronic-based optical response. Nevertheless, in absolute values, the absorption of light by excitons in monolayer TMDs is far below unity, ranging between 2–12% for as-transferred monolayers, and about 20–30% with the aid of a cavity. Similarly, 5–7% absorption have been previously reported for trions. Using thicker TMDs, higher and broadband absorption has been observed, but the thickness was over 20 monolayers.

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complex cavities have also been proposed, but have not been realized to date. Thus, the question remains whether the achievable interaction strength can be pushed further and what would its limit be? Can unitary absorption be experimentally reached by excitons in an atomic thin layer? The answers may play an important role in the understanding of excitonic complexes in TMDs, and the realization of practical 2D material optoelectronic devices.

Here, we demonstrate ultrastrong light-exciton interaction in a W S 2 -based high quality van der Waals heterostructure cavity (VHC), that is, an optical cavity built from van der Waals (VdW) materials, which can be controlled both electrically and optically. While the cavity is quite broadband, the near-unity absorption is attainable owing to four major elements: (1) achievable narrow excitonic line widths, owing to the VdW heterostructure, (2) the ability to carefully balance the interplay between the radiative and nonradiative decay rates, γ r and γ nr, respectively, (3) the enhancement of the vacuum radiative decay rate, γ r,0 via the Purcell factor, shifting the maximum absorption to larger (and more attainable) line widths, and (4) obtaining extremely low dephasing rate, γ d ≪ γ r, γ nr.

We show that this approach yields a large photoexcited excitonic population, with record values of ~95% excitonic absorption, ~41% for singlet and triplet trion states, and even the observation of the next negatively charged trion state with ~28% absorption. In addition, it enables the observation of biexcitons photoluminescence (PL) at ultralow continuous-wave (cw) laser powers down to a few nW, which is 3 orders of magnitude lower than previously reported values. We introduce an analytical approach to describe the light-exciton-cavity interaction, which is based on the semiconductor Bloch equations combined with a quantum transfer matrix method. The model takes into account the contribution of both the exciton radiative and nonradiative decay rates, γ r and γ nr, which have already been shown to affect the high reflection of monolayer TMDs. In addition, we include the existence of a pure dephasing rate, γ d, to account for loss of coherence during the multiple interferences within the cavity. We find that the relation between γ r and γ nr establishes the condition for maximal absorption, and its limit is set by the value of the pure dephasing rate γ d, basically limiting the coherence of the system. Experimentally, we control the nonradiative channels with temperature and the radiative channels via the geometrical parameters of the VHC (Purcell effect). Finally, we demonstrate the existence of a universal absorption law for excitons in 2D systems in this class of devices.

The VHC is composed of a monolayer WS 2 encapsulated by hexagonal-boron-nitride (hBN), with top and bottom thicknesses of 15 and 30 nm, respectively, and transferred on top of a gold back-reflector (Figure 1a). From the optical point of view, this structure forms an asymmetric Fabry-Perot cavity. The simplest case of such a cavity with a thin absorbing layer on top of a dielectric layer and a mirror is known as the Salisbury screen. In this configuration, the thickness of the dielectric is set to a quarter wavelength of the light to achieve constructive interference at the position of the absorbing layer.

**Figure 1.** Light-exciton interaction in a van der Waals heterostructure cavity. (a) The structure of the VHC: a 15 nm hBN/monolayer WS 2 /30 nm hBN heterostructure with optimized hBN thicknesses is transferred on top of a gold back-reflector. The gate voltage V gate controls the doping in the WS 2. (b) Gate-dependent spectral absorption for sample U2, showing ~85% excitonic absorption, ~41% trion absorption, and ~28% of the next negatively charged trion state, at voltages V gate = −0.5 V, V gate = 4.5 V, and V gate = 11 V, respectively. (c) Temperature-dependent excitonic spectral absorption from sample U1, showing the nontrivial behavior of the absorption, and a maximum absorption of ~92% at T = 110 K. (d) Maximum excitonic absorption dependency on the total line width γ T for sample U1 (blue curve) and the temperature-dependent exciton line width (red curve). The black curve shows the measured absorption from sample U3, reaching a maximal value of 95%.
However, these are not sufficient conditions to achieve strong or near-unity abortion, as the amount of absorption is always limited by the physical properties of the thin layer. For example, placing a TMD on top of such a quarter wavelength cavity was predicted to achieve a maximum of \( \sim 33\% \) absorption.\(^{34}\) However, building a TMD-incorporated cavity, as we implemented with the VHC, enables in situ tuning and balancing of the nonradiative and radiative rates of high-quality excitons, which is the key factor in controlling the highest possible absorption. Combining this with careful design of the cavity and integrating it with the 2D-material heterostructure leads to an attainable near-unity absorption of an atomically thin monolayer. Specifically for the VHC, the additional top hBN layer, which is required for the encapsulation, together with the penetration depth into the gold mirror, require an optimization of the final structure (see Supporting Information, SI).

Owing to the well-known inhomogeneity of TMDs,\(^{47}\) we have fabricated several VHC samples: in sample U1, we used for the back-reflector a single crystalline, atomically flat gold flake,\(^{49}\) while in samples U2 and U3, we used a standard evaporated gold film as the back reflector, which in sample U2 has also been used for electrostatic gating. In these structures, the optical transmission (for visible light) is zero, and the absorption can be obtained from \( \frac{1}{\xi} \left( \frac{1}{\xi} \right) \), where \( \xi \) and \( \xi_0 \) are the reflection from the structure with and without the TMD, respectively.

A typical absorption spectra obtained from sample U2 for different gate voltages at \( T = 4 \) K is presented in Figure 1b. An absorption value as high as \( \sim 85\% \) can be seen at an energy of \( E = 2.08 \) eV, corresponding to the WS\(_2\) neutral exciton (X), an absorption value of \( \sim 41\% \) at \( E = 2.041 \) eV, corresponding to both singlet and triplet trions (Tr\(_t/Tr_s\)) (see SI), and an absorption value of \( \sim 28\% \) at \( E = 2.023 \) eV, corresponding to the next charged state of the trio should be marked with two superscript minuses, X\(^{-−}\) (see SI). All energetic positions of the absorption peaks and separations are in agreement with previous reports on WS\(_2\).\(^{6,50,51}\) By changing the gate voltage we control the charge carriers in the WS\(_2\), and thus the relative spectral weights of (X), (Tr\(_t/Tr_s\)) and (X\(^{-−}\)),\(^{36,51}\) that is, the absorption spectrum can be controlled electrically. To the best of our knowledge, these are record absorption values together with the first observation of the (X\(^{-−}\)) peak in an absorption spectra of TMDs. These are a direct result of the strong light-matter interaction provided by the VHC.

To study the ultimate absorption limits, we vary the temperature, which controls the excitonic line width, as presented in Figure 1c and d. While the exciton line width shows a continuous decrease with decreasing temperature (Figure 1d red curve), as known for the ground state exciton in semiconductors\(^{52}\) and TMDs,\(^{42,53}\) the excitonic absorption shows a nonmonotonic temperature dependence. An absorption value of 55% can already be seen at room-temperature, which increases to a maximum value of \( \sim 92\% \) at \( T = 110 \) K, and then decreases rapidly to \( \sim 77\% \) at \( T = 4 \) K (Figure 1d).
To understand the physical origin of this behavior, we developed a theoretical formalism that combines an equation of motion method for the exciton, which is similar to the well-known (interacting) Bloch equations,\textsuperscript{54,55} together with a quantum transfer matrix method (QTMM, see SI). To take into account the pure dephasing, the QTMM treats the electromagnetic fields as operators, rather than classical fields. This provides the relation between the polarization operator and the field in the 2D material. The solution, which describes the optical response of the TMD, leads to an Elliott-type formula appropriate for the 2D material.\textsuperscript{55} This allows us to calculate the absorption, via the reflection operator’s expectation value, taking into account the contributions of both $\gamma_{T,0}$ (and its Purcell enhancement), $\gamma_{nr}$, and $\gamma_d$. The latter is important to consider, not only since the absorption is affected by pure dephasing\textsuperscript{46} but also since a simple classical Gaussian broadening model cannot describe well the experimental behavior (see SI).

Following this approach (see SI), the maximum absorption can then be approximated by

$$A_{\text{max}} = \frac{\gamma_{T,0}}{\gamma_T} \left[ 1 - \xi \left( 1 + 2 \frac{\gamma_d}{\gamma_T} \right) \right]$$

(1)

with

$$\gamma_T = \gamma_{nr} + 2\gamma_d + \xi \gamma_{T,0}$$

(2)

where $\gamma_{nr}$ and $\gamma_d$ are the normalized vacuum radiative decay and dephasing rates, and the coefficients $\xi, \xi_d$, and $\xi$ are (see SI) parameters depending on the geometry, defined by the sizes and dielectric functions composing the different parts of the cavity, with $\xi$ representing the Purcell factor, and $\xi_d = \xi \gamma_{T,0}$ being the renormalized radiative decay rate. As the trion contribution to the susceptibility is not Lorentzian owing to the self-energy term that includes exciton–electron interaction, which stems from the two-particle Green’s function,\textsuperscript{46,57} the same approach cannot be applied for the analysis of the trion response.

To expose the roles of the different decay rates in the absorption behavior, we show in Figure 2a the calculated absorption dependence on the exciton line width, $\gamma_T$, for several dephasing values. For each $\gamma_d$, $\gamma_{T,0}$ is kept constant, as it does not depend on temperature, and we vary $\gamma_{nr}$ (solid lines). For the simple case of negligible $\gamma_d$, we obtain a matching condition for the maximum absorption point: $\gamma_{nr} \approx \gamma_{T,0}$ (see SI), and 100% absorption. This implies that the relation between $\gamma_{nr}$ and $\gamma_{T,0}$ instead of their absolute values, sets the matching condition for maximal absorption, and $\gamma_d$ sets the limit of the absolute achievable absorption when the matching condition is fulfilled. For comparison, the absorption of a suspended monolayer is presented for the same dephasing values (dashed curves), exhibiting the same behavior and showing the known maximal absorption limit of 50% for a thin layer,\textsuperscript{58} but at smaller and less attainable conditions.

The above-discussed decay rates are highly dependent on the cavity design, TMD quality, fabrication-induced interface quality, and can vary locally.\textsuperscript{47} Yet, via a simple representation of eq 1 as a function of $\gamma_T$, the model makes a striking prediction and inescapable universal feature of this class of devices:

$$\frac{A_{\text{max}}}{\xi} = \frac{\gamma_{T,0}}{\gamma_T} \left[ 1 - \xi \left( 1 + 2 \frac{\gamma_d}{\gamma_T} \right) \right]$$

(3)
Equation 3 implies a linear relationship between $A_{\text{max}} T/\gamma^2$ and $1/\gamma T$, provided that $\gamma_d/\gamma_T \ll 1$. Indeed, the universal law is confirmed by the experimental data from different samples and locations presented in Figure 2c. The different samples follow their own straight line, which encodes their different quality, but the generic behavior is the same for all. In addition, $\gamma_d$ of the different samples can be extracted from the (extrapolated) crossing point $1/\gamma_T = 0$. In principle, this universal law should hold for any 2D excitonic system where $\gamma_d \ll 1$.

Another important outcome of this analysis is the ability to compare different excitonic properties via their absorption response. Figure 2d shows the extracted temperature-dependent $\gamma_w$ and $\gamma_d$ (see SI) and their phenomenological fit. These correlate directly with their absorption behavior, indicating that higher dephasing leads to lower absorption.

In the same manner, we have designed the VHC for maximal interaction strength, it can be designed to any intermediate value, and even to completely turn off the interaction between the light and the TMD. To demonstrate this, we fabricated on the same device two different cavities, one for which the interaction is optimized (denoted as “on”) and one for which the interaction is minimized (denoted as “off”). This is done by adding another hBN flake below a part of the heterostructure. The spatial distribution of the absorption for this device at $T = 300\,\text{K}$ is presented in Figure 3a (see SI for exemplary spectral curves) and its extracted exciton line width distribution in Figure 3b. The spatial correlation between the absorption and line width can be directly observed in the “on” areas of the two figures, that is, lower line width correlates to higher absorption, which is indeed the case for $T = 300\,\text{K}$ (see Figure 1d). The source of the spatial distribution comes from the inhomogeneity of the sample, a well-known issue in TMDs.47 Furthermore, the extracted $\gamma$ from the two cavities yields 2.2 meV and 70 $\mu\text{eV}$ in the “on” and “off” regions, respectively. This is in agreement with the Purcell effect line width modulation obtained by refs 59–62.

The ability to achieve near-unity excitonic absorption implies that a large photoexcited excitonic density can be obtained, while maintaining low excitation power. It was already shown that TMDs are highly affected by the excitation power, resulting in either heating effects that changes the excitonic properties temporarily,63 or permanent effects that completely alter the material’s response, such as optical doping and environmental surface interactions.64 Yet, this challenging and desirable high exciton density plays a major role in several physical phenomena, such as Bosonic condensation, phase transitions,65 and biexcitons emission,\textsuperscript{5} for example. The formation of biexcitons is directly related to high excitonic density and thus also to the possible appearance of the biexciton peak in the PL spectrum.\textsuperscript{5} The predicted intensity relation between the exciton and biexciton emission is a power law, $I_{XX} = I_{XXr}^\alpha$, with $\alpha$ ranging between 1.2–1.9 due to lack of thermal equilibrium.\textsuperscript{5} Thus, the latter can be used to probe the efficiency of exciton photogeneration in the VHC.

Figure 4a shows the PL spectra, normalized to the excited emission intensity ($I_{Xr}$), obtained from the VHC for different cw excitation powers. Several peaks can be observed and are marked as X, exciton at $E = 2.071\,\text{eV}$; $T_X/T_{XX}$ singlet/triplet trions\textsuperscript{50} at 2.034 eV/2.04 eV, respectively; and XX, biexciton at $E = 2.018\,\text{eV}$. Remarkably, the XX emission peak can be observed at excitations powers down to few $\text{nW}$. This excitation power is 2–3 orders of magnitude smaller than the lowest previously reported for biexcitons.\textsuperscript{44} Actually, for excitation powers above 30 $\mu\text{W}$, the XX emission intensity is so high that it saturates our detector. Figure 4b shows the power law analysis of the $I_{XX}$ and $I_X$ peaks, and the obtained $\alpha = 1.62$ for $I_{XX}$ confirms the identity of the biexciton. From the maximal used power of 400 $\mu\text{W}$, we can calculate an excitonic density of $\sim 2.5 \times 10^{11}\,\text{cm}^{-2}$. This exciton density is 2 orders of magnitude larger than what was achieved in previous works using cw excitation, and is similar to densities obtained using ultrafast laser pulses.\textsuperscript{29} While it is possible that at high cw excitation powers heating effects may kick in, using pulsed excitation (as in ref 29 for example) with our VHC, one can obtain an excitonic density of $\sim 5 \times 10^{13}\,\text{cm}^{-2}$ with little heating, as compared to $\sim 6.4 \times 10^{12}\,\text{cm}^{-2}$.

In conclusion, the new type of optical cavity presented here, which is based solely on vdW heterostructures, is designed for ultrastrong light-exciton interaction, and unveils a universal absorption law for excitons in 2D systems. This enhanced light-exciton interaction can act as a platform for probing quantum nonlinear dynamics of excitons,\textsuperscript{46,48} and their state of matter,\textsuperscript{65} and paves the way to efficient optoelectronic devices, such as detectors, modulators, and optically pumped light emitting devices, based on monolayer semiconductors. Specifically, this high absorption is important for TMD-based photodetectors, as the internal photodetection efficiency of TMDs has been demonstrated to be high,\textsuperscript{66} but the external efficiency has always been limited by the low absorption efficiency.
ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/10.1021/acs.nanolett.0c00492.

Theoretical modeling, fitting procedure, identification of excitonic complexes, structure optimization, and experimental and fabrication methods (PDF)

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Notes

The authors declare no competing financial interest.

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