THE SUPER FROBENIUS-SCHUR INDICATOR AND
FINITE GROUP GAUGE THEORIES ON PIN⁻ SURFACES

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ABSTRACT. It is well-known that the value of the Frobenius-Schur indicator $|G|^{-1} \sum_{g \in G} \chi(g^2) = \pm 1$ of a real irreducible representation of a finite group $G$ determines which of the two types of real representations it belongs to, i.e. whether it is strictly real or quaternionic. We study the extension to the case when a homomorphism $\varphi : G \to \mathbb{Z}/2\mathbb{Z}$ gives the group algebra $\mathbb{C}[G]$ the structure of a superalgebra. Namely, we construct a super version of the Frobenius-Schur indicator whose value for a real irreducible super representation is an eighth root of unity, distinguishing which of the eight types of irreducible real super representations described in [Wal64] it belongs to. We also discuss its significance in the context of two-dimensional finite-group gauge theories on pin⁻ surfaces.

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1. INTRODUCTION

It is a classic result of Frobenius and Schur [FS06] that for a finite group $G$ and its irreducible representation $\rho$ over $\mathbb{C}$ whose character we denote by $\chi$, the Frobenius-Schur indicator defined by

$$S(\rho) := \frac{1}{|G|} \sum_{g \in G} \chi(g^2)$$

(1.1)

takes values in $\{\pm 1, 0\}$, depending on whether $\rho$ is strictly real, complex or quaternionic. A slight reformulation can be given in terms of the simple summand $A_\rho$ of the group algebra $\mathbb{C}[G]$. When $\rho \simeq \bar{\rho}$, $A_\rho$ has a conjugate-linear involution $\ast$, whose fixed point is a central simple algebra over $\mathbb{R}$. From Wedderburn’s theorem, this is a matrix algebra over a division algebra $D$ over $\mathbb{R}$ whose center is $\mathbb{R}$, which is either $\mathbb{R}$ or $\mathbb{H}$. These algebras $\mathbb{R}$, $\mathbb{H}$ form the Brauer group.
In this paper we generalize these statements to the group superalgebras. Let us first recall the Brauer-Wall group $BW(F)$ of a field $F$ [Wal64, Del99]. A division superalgebra $D = D_0 \oplus D_1$ is a superalgebra such that every nonzero element in $D_0$ or $D_1$ is invertible. Consider a simple superalgebra $A = A_0 \oplus A_1$ over $F$, with its unique irreducible supermodule $\rho$. The super version of Schur’s lemma says that $D_\rho := \text{Hom}_A(\rho, \rho)$ is a division superalgebra over $F$. Simple superalgebras $A$ such that the center of the even part $(D_\rho)_0$ is $F$ itself are called central simple, and central simple superalgebras under the equivalence relation $A \sim A' \iff D_\rho \simeq D_{A'}$ form a finite group $BW(F)$ under tensor product. This is the Brauer-Wall group of $F$. It was determined in [Wal64] that $BW(\mathbb{C}) = \mathbb{Z}/2\mathbb{Z}$ and $BW(\mathbb{R}) = \mathbb{Z}/8\mathbb{Z}$. The Clifford algebras over $\mathbb{C}$ and $\mathbb{R}$, with the corresponding periodicity 2 and 8, cover all ten cases.

Let us now specialize to the case of superalgebras associated to a group, and formulate our main theorem. Let $G$ be a finite group, $\varphi : G \to \mathbb{Z}/2\mathbb{Z}$ a homomorphism or equivalently a one-cocycle, and $\alpha : G \times G \to \mathbb{Z}/2\mathbb{Z}$ a two-cocycle. $\varphi$ gives a decomposition of $G$ as $G = G_0 \sqcup G_1$. We define $\mathbb{R}[G]_{\varphi, \alpha}$ to be the superalgebra over $\mathbb{R}$ generated by elements $e_g$ for $g \in G$ such that $e_g$ is even or odd depending on $\varphi(g) = 0$ or 1, and $e_g e_h = (-1)^{\alpha(g,h)} e_{gh}$. We denote its complexification by $A := \mathbb{C}[G]_{\varphi, \alpha}$. This algebra is equipped with a conjugate linear involution $\alpha \mapsto \alpha^\ast$, inherited from the complex conjugation of $\mathbb{C}$. These superalgebras are semisimple, as can be seen e.g. by using the unitary trick. For a supermodule $\rho$ defined on $V = V_0 \oplus V_1$, we define its character as $\chi(g) := \text{tr}_V \rho(e_g)$.

Let $(\rho, V)$ be an irreducible supermodule of $A = \mathbb{C}[G]_{\varphi, \alpha}$. The corresponding subsuperalgebra $A_\rho$ of $A$ is central simple, and we denote its class $[A_\rho] \in BW(\mathbb{C})$ by $q(\rho) = 0, 1$. It is known that $\rho$ is reducible as an ungraded module if $q(\rho) = 0$, and $\rho|_{V_0}$ is irreducible as a module of $\mathbb{C}[G_0]_\alpha$ if $q(\rho) = 1$.

We denote by $\bar{\rho}$ the complex conjugate supermodule of $\rho$. When $\bar{\rho} \simeq \rho$, $\rho$ is called real. Otherwise $\rho$ is called complex. Real irreducible supermodules of $\mathbb{C}[G]_{\varphi, \alpha}$ can be classified by $BW(\mathbb{R}) = \mathbb{Z}/8\mathbb{Z}$, just as irreducible supermodules of $\mathbb{R}[G]_{\varphi, \alpha}$ can be.

We show the following:

**Theorem 1.2.** The super Frobenius-Schur indicator of an irreducible supermodule $\rho$, defined by

$$S(\rho) := \frac{1}{\sqrt{2^{|G|}}} \sum_{g \in G} \sqrt{-1}^{\varphi(g)} (-1)^{\alpha(g,g)} \chi(g^2),$$

is either zero or an eighth root of unity. It is zero if $\rho$ is complex and non-zero if $\rho \simeq \bar{\rho}$. In the latter case, its value $S(\rho) \in \{x^8 = 1 \mid x \in \mathbb{C}\}$ determines to which element of $BW(\mathbb{R}) \simeq \mathbb{Z}/8\mathbb{Z}$ the supermodule belongs.

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1Various other generalizations of the Frobenius-Schur indicator were studied in the literature. Firstly, we can consider generalized Frobenius-Schur indicators for finite groups, see e.g. [BG04] and [GI18, Sec. 11]. Secondly, we can extend the concept of the Frobenius-Schur indicator to fusion categories, see e.g. [FS01, Sec. 3], which is further extended to fermionic fusion categories in e.g. [BWHV17, Sec. 3.3]. This last generalization should agree with ours when specialized to fermionic fusion categories arising from our superalgebra. It would be interesting to check this.
The rest of the paper consists of two sections. In Section 2, we provide a proof of the main theorem. In Section 3, we describe how the super Frobenius-Schur indicator appears in the context of two-dimensional finite-group gauge theories on pin$^-$ surfaces.

2. Proof of the main theorem

We start by quoting the structure of central simple complex or real superalgebras $A = A_0 \oplus A_1$ and the corresponding Brauer-Wall groups $BW(\mathbb{C}) = \mathbb{Z}/2\mathbb{Z}$ or $BW(\mathbb{R}) = \mathbb{Z}/8\mathbb{Z}$, as described in [Wal64, Del99]. For simplicity of the presentation, we assume that any superalgebra $A$ contains a regular element in $A_1$.

Over $\mathbb{C}$, any central simple superalgebra $A = A_0 \oplus A_1$ is either of the form

\begin{equation}
M(n) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \right\}
\end{equation}

or

\begin{equation}
Q(n) = \left\{ \begin{pmatrix} a & b \\ b & a \end{pmatrix} \right\}
\end{equation}

where $a, b, c, d$ are $n \times n$ matrices. $a$ and $d$ provide even elements and $b$ and $c$ provide odd elements. $M(n)$ belongs to $0 \in BW(\mathbb{C})$ and $Q(n)$ belongs to $1 \in BW(\mathbb{C})$. We denote the type in $BW(\mathbb{C})$ by $q \in BW(\mathbb{C})$. We note that $Q(1)$ is the Clifford algebra $Cl(1)$.

More abstractly, any central simple superalgebra over $\mathbb{C}$ either admits a nonzero graded central element $u \in A_0$ or a nonzero central element $u \in A_1$, depending on whether $q = 0$ or $1$. Here, an element $x$ is called graded central if $yx = xy$ for all $y \in A_0$ and $yx = -xy$ for all $y \in A_1$. The special element $u$ is determined up to a scalar multiplication. We note that $A$ or $A_0$ is simple as an ungraded algebra when $q = 0$ or $q = 1$, respectively.

By direct computation, one finds that

\begin{equation}
M(m) \hat{\otimes} M(n) = M(mn), \quad M(m) \hat{\otimes} Q(n) = Q(mn), \\
Q(m) \hat{\otimes} M(n) = Q(mn), \quad Q(m) \hat{\otimes} Q(n) = M(mn) \oplus M(mn)
\end{equation}

where $\hat{\otimes}$ denotes the super tensor product of two superalgebras.

An irreducible supermodule $(\rho, V)$ of the semisimple superalgebra $A = \mathbb{C}[G]_{\varphi, \alpha}$, determines a simple summand $A_\rho$ of $A$. This allows us to associate an element $q \in BW(\mathbb{C})$ to the supermodule $\rho$. We note that $\rho$ is irreducible as an ungraded module when $q = 0$, and $\rho \mid_{V_0}$ is irreducible as an ungraded module of $\mathbb{C}[G_0]_\alpha$ when $q = 1$. We recall that $G_0$ is the inverse image of $0 \in \mathbb{Z}/2\mathbb{Z}$ under $\varphi : G \to \mathbb{Z}/2\mathbb{Z}$.

Let us take another finite group $H$, a homomorphism $\varphi' : H \to \mathbb{Z}/2\mathbb{Z}$, and a $\mathbb{Z}/2\mathbb{Z}$-valued two-cocycle $\alpha'$ of $H$. Let us consider a supermodule $\rho$ of $\mathbb{C}[G]_{\varphi, \alpha}$ and a supermodule $\sigma$ of $\mathbb{C}[H]_{\varphi', \alpha'}$.

**Lemma 2.4.** $\rho \hat{\otimes} \sigma$ is a supermodule of $\mathbb{C}[G \times H]_{\varphi, \alpha \varphi'}$ where

\begin{equation}
\varphi^0 = \varphi + \varphi', \quad \alpha^0 = \alpha + \alpha' + \varphi \varphi'.
\end{equation}

Here, on the right hand side, $\varphi, \varphi', \alpha, \alpha'$ are appropriately pulled back from $G$ and $H$ to $G \times H$. Then, to take the product of $\varphi$ and $\varphi'$, we regard them as one-cocycles so that the result is a two-cocycle, i.e., $(\varphi \varphi')(g, h) := \varphi(g) \varphi'(h)$. 

Proof. Direct computation. □

Let us assume that $\rho$ and $\sigma$ are irreducible. According to (2.3), $\hat{\rho} \otimes \sigma$ is irreducible or a direct sum of two isomorphic irreducible supermodules. We denote by $\rho \sigma$ the isomorphism class of the irreducible supermodule(s) contained in $\rho \hat{\otimes} \sigma$.

**Lemma 2.5.** The super Frobenius-Schur indicator is multiplicative in the sense that 
\[ S(\rho \sigma) = S(\rho)S(\sigma). \]

Proof. Direct computation. □

Let us further assume that $(\rho, V)$ is real, i.e. $\rho \simeq \bar{\rho}$. In this case the irreducible representations are classified into eight types in the following manner:

- We first ask whether the corresponding $q \in \text{BW}(\mathbb{C})$ is 0 or 1. This gives the first two-fold division.
- We then replace the special element $u$ by $cu$ for some $c \in \mathbb{C}$ so that we have $u = u^*$ and $u^2 = \pm 1$. The sign gives the second two-fold division.
- Finally, when $q = 0$, $\rho$ is real irreducible as an ungraded module of $\mathbb{C}[G]_{\varphi, \alpha}$, and we ask whether $\rho$ is strictly real or quaternionic. Similarly, when $q = 1$, $\rho|_{V_0}$ is real irreducible as an ungraded module of $\mathbb{C}[G_0]_{\alpha}$, and we ask whether $\rho|_{V_0}$ is strictly real or quaternionic. This gives the third two-fold division.

We note that when $\rho \simeq \bar{\rho}$, the simple summand $A_\rho$ of $\mathbb{C}[G]_{\varphi, \alpha}$ has a conjugate-linear automorphism $\ast$, whose fixed points form a central simple superalgebra over $\mathbb{R}$. We recall that central simple superalgebras over $\mathbb{R}$ can be classified in terms of $\text{BW}(\mathbb{R}) = \mathbb{Z}/8\mathbb{Z}$, whose structure was stated exactly as above [Wal64, p.195]. We summarize the structure in the following table:

| $\text{BW}(\mathbb{R})$ | $q$ | $u^2$ | $\rho$ | $\text{BW}(\mathbb{R})$ | $q$ | $u^2$ | $\rho|_{V_0}$ |
|---|---|---|---|---|---|---|---|
| 0 | 0 | +1 | $\mathbb{R}$ | 1 | 1 | +1 | $\mathbb{R}$ |
| 2 | 0 | −1 | $\mathbb{R}$ | 3 | 1 | −1 | $\mathbb{H}$ |
| 4 | 0 | +1 | $\mathbb{H}$ | 5 | 1 | +1 | $\mathbb{H}$ |
| 6 | 0 | −1 | $\mathbb{H}$ | 7 | 1 | −1 | $\mathbb{R}$ |

(2.6)

Let us now consider some examples.

**Lemma 2.7.** Let $G = (\mathbb{Z}/2\mathbb{Z})^n$. There is a homomorphism $\varphi^{(n)} : G \to \mathbb{Z}/2\mathbb{Z}$ and a $\mathbb{Z}/2\mathbb{Z}$-valued two-cocycle $\alpha^{(n)}$ on $G$ such that $\mathbb{C}[G]_{\varphi^{(n)}, \alpha^{(n)}}$ is the Clifford algebra $\text{Cl}(n)$.

Proof. Take $G^o = \mathbb{Z}/2\mathbb{Z}$, $\varphi^o : \mathbb{Z}/2\mathbb{Z} \to \mathbb{Z}/2\mathbb{Z}$ is the identity, and $\alpha^o \equiv 0$. We have $A^o := \mathbb{C}[G^o]_{\varphi^o, \alpha^o} = \mathbb{C}[u]$, where $u$ is odd and $u^2 = 1$. This equals the Clifford algebra $\text{Cl}(1)$. We note that $(A^o)^{\otimes n}$ is the Clifford algebra $\text{Cl}(n)$. From the lemma 2.4 above, this is a twisted group superalgebra associated to $G = (\mathbb{Z}/2\mathbb{Z})^n$ with a specific choice of $\varphi^o$ and $\alpha^o$. □

By comparing the explicit description of $A^o = \text{Cl}(1)$ and the table 2.6, we find that $A^o = \text{Cl}(1)$ corresponds to $1 \in \text{BW}(\mathbb{R})$. Therefore $\text{Cl}(n)$ corresponds to $n \in \text{BW}(\mathbb{R})$. It has a unique irreducible supermodule.
Lemma 2.8. For the irreducible supermodule $\rho$ of $\Cl(n) = \mathbb{C}[(\mathbb{Z}/2\mathbb{Z})^n]_{\varphi(n),\alpha(n)}$ defined above, the super Frobenius-Schur indicator is given by

$$S(\rho) = e^{n \times 2\pi \sqrt{-1}/8}.$$ 

Proof. For $n = 1$ this follows from a direct computation. Then the statement for $n > 1$ follows from a repeated application of the lemma 2.5.

We note that the summand $\sqrt{-1}^{e(g)}(-1)^{\alpha(g,h)} : (\mathbb{Z}/2\mathbb{Z})^n \to \{\pm 1, \pm \sqrt{-1}\}$ in the definition of the indicator in this case is a $\mathbb{Z}/4\mathbb{Z}$-valued quadratic refinement of the standard non-degenerate pairing on $(\mathbb{Z}/2\mathbb{Z})^n$, and that the super Frobenius-Schur indicator is its Arf-Brown invariant itself.

We also need the following lemma:

Lemma 2.9. Let $\mathbb{C}[G]_\alpha$ be the twisted group algebra generated by $e_g$ for $g \in G$ satisfying $e_ge_h = (-1)^{\alpha(g,h)}e_{g,h}$, where $\alpha(g,h) \in \{0,1\}$ is a two-cocycle. Its irreducible representation $\rho$ is complex, strictly real, or quaternionic if and only if its Frobenius-Schur indicator

$$S(\rho) := \frac{1}{|G|} \sum_e (-1)^{\alpha(g,h)} \chi(g^2)$$

is $0$, $+1$ or $-1$, respectively.

Proof. We repeat the standard proof which studies whether $\rho \otimes \rho$ contains the identity representation of $G$, and if so, whether it is in $\text{Sym}^2 \rho$ or in $\wedge^2 \rho$. We find that

$$\frac{1}{|G|} \sum_g \text{tr}(e_ge_g) = 0, +1, -1$$

depending on the three cases. Finally we use $e_ge_g = (-1)^{\alpha(g,h)} \chi(g^2)$.

We can now begin the

Proof of the main theorem. The homomorphism $\varphi : G \to \mathbb{Z}/2\mathbb{Z}$ gives the decomposition $G = G_0 \sqcup G_1$. When $G_1$ is empty, the theorem reduces to the non-super case, our Lemma 2.9. When $\alpha$ is trivial it further reduces to the original theorem of Frobenius-Schur [FS06]. We assume $G_1$ is non-empty in the following. We then have $|G| = 2|G_0|$.

We saw in lemma (2.5) that the super Frobenius-Schur indicator is multiplicative. We also saw in lemma (2.8) that the super Frobenius-Schur indicator of the irreducible supermodule $\rho^\circ$ of $\Cl(1) = \mathbb{C}[G^\circ]_{\varphi^\circ,\alpha^\circ}$ is $e^{2\pi \sqrt{-1}/8}$. Therefore, repeating $n$ times the procedure of tensoring by $\rho^\circ$ and taking the irreducible component multiplies the indicator by $e^{2\pi \sqrt{-1}n/8}$ and adds $n$ mod 8 to the class in $\text{BW}(\mathbb{R})$. This means that it suffices to prove our statement in two cases, namely when $\rho$ is complex and corresponds to $0 \in \text{BW}(\mathbb{C})$, or when $\rho$ is real and corresponds to $0 \in \text{BW}(\mathbb{R})$. Below, we assume that either of the two is satisfied.

We note that $S(\rho)$ can be rewritten as

$$(2.10) \quad S(\rho) = \frac{1}{\sqrt{2^n}} \left[ \frac{1}{|G_0|} \sum_{g \in G_0} (-1)^{\alpha(g,h)} \chi(g^2) + \frac{\sqrt{-1}}{|G|} \left( \sum_{g \in G} - \sum_{g \in G_0} \right) (-1)^{\alpha(g,h)} \chi(g^2) \right]$$

This means that $S(\rho)$ can be computed if the ordinary Frobenius-Schur indicators of $\rho$ as representations of $G$ and $G_0$ can be computed.
Denote $V = V_0 \oplus V_1$ the graded vector space on which $\rho$ is defined. Due to our assumption, $V$ corresponds to $0 \in \text{BW}(\mathbb{C})$. Therefore $V$ is irreducible as an ungraded module of $\mathbb{C}[G]_{\varphi, \alpha}$. Since $G_1$ is nonempty, $\dim V = 2 \dim V_0$ and $G_0$ is an index-2 subgroup of $G$. Therefore, $V_0$ is an irreducible module of $\mathbb{C}[G_0]_{\alpha}$ and $V$ is the induced representation. Therefore, $V$ is real if $V_0$ is real, and $V$ is complex if $V_0$ is complex. Therefore, when $\rho$ is complex, both $V$ and $V_0$ are irreducible representations of $\mathbb{C}[G]$.

□

There is also an alternative proof, based on an earlier work by Roderick Gow [Gow79], which we briefly indicate below\(^2\). Given an index-2 subgroup $G_0$ of $G$ and an irreducible representation $\rho_0$, Gow considered an analogue of the Frobenius-Schur indicator

\begin{equation}
\eta(\rho_0) := \frac{1}{|G_0|} \sum_{g \in G \setminus G_0} \chi(g^2),
\end{equation}

which we call the Gow indicator. In [Gow79, Lemma 2.1], it was proved that $\eta(\rho_0) = +1, 0, -1$, whose value depends on $S(\rho_0)$ and various properties of the $G$ representation $\text{Ind}_{G_0}^G(\rho_0)$ induced from $\rho_0$ of $G_0$; see also [KM90, Sec. 2]. The analysis there can be readily generalized to the case when the $\mathbb{Z}/2\mathbb{Z}$-valued $2$-cocycle $\alpha$ is nontrivial.

For us, we have $\rho_0 = \rho|_{V_0}$ and $\rho = \text{Ind}_{G_0}^G(\rho_0)$. Then $q = 0, 1$ depending on whether $\rho$ is irreducible or not as a $G$ representation. We can then straightforwardly show that

\begin{equation}
S(\rho) = \frac{1}{\sqrt{2^q}} \left( S(\rho_0) + \sqrt{-1} \eta(\rho_0) \right).
\end{equation}

This allows us to determine our super Frobenius-Schur indicator from the description of the Gow indicator in [Gow79, Lemma 2.1], by carefully comparing its proof and our table (2.6). This gives an alternative proof.

Another related point is that the element of the Brauer-Wall group associated to an irreducible representation $\rho$ of an index-2 subgroup $G_0$ of $G$ was studied in detail in [Tur92]. The discussion there was for an arbitrary base field, and no explicit formula was given for the base field $\mathbb{R}$ as in our main theorem.

3. Relation to Two-Dimensional Finite Group Gauge Theories on Pin\(^-\) Surfaces

According to physicists, quantum gauge theories are ‘defined’ in terms of path integrals. We start from a $d$-dimensional manifold $M_d$, and consider the moduli space $\mathcal{M}$ of connections $A$ on principal $G$-bundles on $M_d$. We pick a function $S : \mathcal{M} \to \mathbb{C}/\mathbb{Z}$ called the action. The fundamental object of physicists’ studies is the partition function

\begin{equation}
Z_{G,S}(M_d) := \int_{\mathcal{M}} e^{2\pi \sqrt{-1} S} d\mu
\end{equation}

where $d\mu$ is a suitable measure on $\mathcal{M}$. Of course all this is terribly ill-defined in most of the physically relevant cases when $G$ is a compact Lie group, since $\mathcal{M}$ is heavily infinite-dimensional.

\(^2\)The following three paragraphs were added in v2. The authors thank J. Murray for the information.
Finite group gauge theories, however, admits a path-integral definition without any problems, since the space $M$ over which we need to integrate is a finite set, which is the space of homomorphisms $f : \pi_1(M_d) \to G$ up to conjugation. The proper measure $d\mu$ to be used was determined e.g. in [FQ91], and the partition function is simply

\begin{equation}
Z_{G,\mathbb{S}}(M_d) := \frac{1}{|G|} \sum_{f : \pi_1(M_d) \to G} e^{2\pi\sqrt{-1}\mathbb{T}(f)}
\end{equation}

for a closed and connected manifold $M_d$. Here we only described the partition function, but they can be promoted to topological field theories in the sense of Atiyah-Segal, and further to extended topological field theories.

When the space $M = \Sigma$ is two-dimensional and the action $\mathbb{S}$ is trivial, this ‘path integral’ was considered independently of physics in [Med78], which says the following. Let $\Sigma$ be a two-dimensional oriented surface of genus $\gamma$ which is closed and connected. We then have

\begin{equation}
\frac{1}{|G|} \sum_{f : \pi_1(\Sigma) \to G} 1 = \sum_{\rho \in \text{Irr}(G)} \left( \frac{|G|}{\dim \rho} \right)^{-e(\Sigma)}
\end{equation}

where the sum on the left hand side is over homomorphisms $f$ from $\pi_1(\Sigma)$ to $G$, the sum on the right hand side is over isomorphism classes of irreducible representations of $G$, and $e(\Sigma) = 2 - 2\gamma$ is the Euler number of the surface. When $\gamma = 0$ this reduces to the classic formula $|G| = \sum_{\rho \in \text{Irr}(G)} (\dim \rho)^2$.

For the detailed history of these formulas and very elementary proofs, see [Sny07]. This formula is so classic that it appears as exercises of textbooks of finite group theory, e.g. [Ser16].

The formula (3.3) can be further generalized in many directions. One example starts by picking a degree-2 cohomology class $\alpha \in H^2(BG, \mathbb{R}/\mathbb{Z})$. A homomorphism $f : \pi_1(\Sigma) \to G$ determines a classifying map $f : \Sigma \to BG$ of the corresponding $G$-bundle, which we denoted by the same letter. We now take the pull-back $f^*(\alpha)$ of $\alpha$ and integrate it against the fundamental class, the result of which we denote by $\int_{\Sigma} f^*(\alpha)$. We use it as the action $\mathbb{S}$ and the generalized formula is now given by

\begin{equation}
\frac{1}{|G|} \sum_{f : \pi_1(\Sigma) \to G} e^{2\pi\sqrt{-1}\int_{\Sigma} f^*(\alpha)} = \sum_{\rho \in \text{Irr}_\alpha(G)} \left( \frac{|G|}{\dim \rho} \right)^{-e(\Sigma)},
\end{equation}

where $\text{Irr}_\alpha(G)$ is now the set of isomorphism classes of irreducible projective representations of $G$ whose degree-2 cocycle is specified by $\alpha$. The use of the pull-back of a cohomology class of the classifying space as the action goes back to [DW90]. This particular formula, appropriately generalized to compact Lie group $G$, first appeared in [Wit91] to the authors’ knowledge. For early rigorous mathematical exposition, see [FQ91].

The simplest way to allow $\Sigma$ to be non-orientable is to take $\alpha \in H^2(BG, \mathbb{Z}/2\mathbb{Z})$ instead, since any surface $\Sigma$, non-orientable or otherwise, has a fundamental class $[\Sigma] \in H_2(\Sigma, \mathbb{Z}/2\mathbb{Z})$ and we can pair $f^*(\alpha)$ with it. The formula now involves the Frobenius-Schur indicator discussed in
Lemma 2.9,

\[
\frac{1}{|G|} \sum_{f: \pi_1(\Sigma) \to G} e^{2\pi \sqrt{-1} I f_* f^*(\alpha)} = \sum_{\rho \in \text{Irr}(G)} S(\rho) o(\Sigma) \left( \frac{|G|}{\dim \rho} \right)^{-e(\Sigma)}
\]

where \( o(\Sigma) = 0, 1 \) is the class \([\Sigma] \in \Omega_{d=2}^{\text{oriented}} = \mathbb{Z}/2\mathbb{Z} \) and equals \( o(\Sigma) = e(\Sigma) \mod 2 \). This formula also already appeared in [Wit91]; for a rigorous mathematical exposition, see [Tur07].

We also note that for \( \Sigma = \mathbb{R}P^1 \) and for trivial \( \alpha \), the formula reduces to

\[
\sum_{g^2 = 1} 1 = \sum_{\rho \in \text{Irr}(G)} S(\rho) \dim \rho,
\]

a formula which appears already in the very paper of Frobenius and Schur [FS06].

More recently, it has become of interest to consider finite group gauge theories on spacetimes with spin structures and variants such as pin\(^{\pm}\) structures. The motivation came both from developments internal to mathematical physics, see e.g. [BT13, NR14] and also from influences from theoretical condensed matter physics, see e.g. [KTTW14, GK15, WOP+18, GOP+18].

Now, what type of actions \( S \) should we consider? As was first emphasized in [FM04], natural choices of the topological part of the action should be given by invertible quantum field theories; for details of invertible quantum field theories, the readers are referred to excellent lecture notes by Freed [Fre19]. For a finite group gauge theory, they are specified by bordism invariants

\[
\phi \in \text{Hom}(\Omega^\text{structure}_d(BG), \mathbb{R}/\mathbb{Z})
\]

where ‘structure’ in the equation stands for ‘unoriented’, ‘oriented’, ‘spin’ etc., appropriate for the purpose, and \( d \) is the space-time dimension. Therefore, we are led to consider the sum

\[
\frac{1}{|G|} \sum_{f: \pi_1(M_d) \to G} e^{2\pi \sqrt{-1} I \phi([f: M_d \to BG])},
\]

where \( M_d \) is a \( d \)-dimensional closed and connected manifold equipped with the structure of your choice.

For oriented surfaces with a \( G \) bundle, we have

\[
\text{Hom}(\Omega^\text{oriented}_2(BG), \mathbb{R}/\mathbb{Z}) = H^2(BG, \mathbb{R}/\mathbb{Z}),
\]

corresponding to the Dijkgraaf-Witten theories we recalled above. For oriented spin surfaces with a \( G \) bundle, we have an equality as sets

\[
\text{Hom}(\Omega^\text{spin}_2(BG), \mathbb{R}/\mathbb{Z}) = H^1(BG, \mathbb{Z}/2\mathbb{Z}) \times H^2(BG, \mathbb{R}/\mathbb{Z}),
\]

where the group structure on the right hand side is given by

\[
(\varphi, \alpha) + (\varphi', \alpha') = (\varphi + \varphi', \alpha + \alpha' + \iota(\varphi \varphi')).
\]

Here \( \iota \) is the homomorphism \( \mathbb{Z}/2\mathbb{Z} \to \mathbb{R}/\mathbb{Z} \) obtained by sending 1 to 1/2, see [BM16, BM18a] where \( d = 2, 3, 4 \) were treated in parallel.

An element \( \phi = (\varphi, \alpha) \in \text{Hom}(\Omega^\text{spin}_2(BG), \mathbb{R}/\mathbb{Z}) \) is paired with a \( G \)-bundle \( f: \Sigma \to BG \) on a spin surface in the following manner. We first recall [Ati71, Joh80] that a spin structure on a
spin surface can be identified with the quadratic refinement \( Q : H^1(\Sigma, \mathbb{Z}/2\mathbb{Z}) \to \mathbb{Z}/2\mathbb{Z} \), i.e. those functions which satisfy
\[
Q(a + b) - Q(a) - Q(b) = \int_\Sigma a \cup b \tag{3.12}
\]
and \( Q(0) = 0 \). Then we evaluate \( \phi \) by the formula
\[
e^{2\pi \sqrt{-1} \phi([f: \Sigma \to BG])} := (-1)^{Q(f^*(\varphi))} e^{2\pi \sqrt{-1} f^*(\alpha)}. \tag{3.13}
\]
The term \( \iota(\varphi \varphi') \) in (3.11) follows by using the last two formulas.

The generalization of the formula (3.3) is given by
\[
\frac{1}{|G|} \sum_{f: \pi_1(\Sigma) \to G} (-1)^{Q(f^*(\varphi))} e^{2\pi \sqrt{-1} f^*\alpha} = \sum_{\rho \in \text{Irr}(\mathbb{C}[\varphi, \alpha])} (-1)^{\text{Arf}(\Sigma) q(\rho)} \left( \frac{|G|}{\text{qdim } \rho} \right)^{-e(\Sigma)} \tag{3.14}
\]
where the sum on the right hand side is now over the irreducible supermodules \( \rho \) of \( \mathbb{C}[G]_{\varphi, \alpha} \) discussed above. \( q(\rho) \in BW(\mathbb{C}) = \{0, 1\} \) is the type of the supermodule \( \rho \), \( \text{qdim } \rho \) = \( \dim(\rho)/\sqrt{q(\rho)} \), and the Arf invariant \( \text{Arf}(\Sigma) = \mathbb{Z}/2\mathbb{Z} \) is the class \( [\Sigma] \in \Omega^\text{spin}_2 \) = \{0, 1\} of the spin surface in the bordism group. This formula was derived in [Gun12]; the presentation there was mainly for the case \( G = (\mathbb{Z}/2\mathbb{Z})^n \rtimes S_n \) with a specific choice of \( \varphi \) and \( \alpha \), for which \( \mathbb{C}[G]_{\varphi, \alpha} \) is known as the Sergeev algebra, but the discussion there easily generalizes.

Our super Frobenius-Schur indicator plays a role in the case of unoriented surfaces \( \Sigma \) equipped with a \( \text{pin}^- \) structure. In this case we have an equality of sets [BM18b]
\[
\text{Hom}(\Omega^\text{pin}_2(BG), \mathbb{R}/\mathbb{Z}) = H^1(BG, \mathbb{Z}/2\mathbb{Z}) \times H^2(BG, \mathbb{Z}/2\mathbb{Z}) \tag{3.15}
\]
where the addition formula on the right hand side is
\[
(\varphi, \alpha) + (\varphi', \alpha') = (\varphi + \varphi', \alpha + \alpha' + \varphi \varphi'). \tag{3.16}
\]
Note that this has the same form as the formula we saw in Lemma 2.4.

To evaluate \( \phi = (\varphi, \alpha) \in \text{Hom}(\Omega^\text{pin}_2(BG), \mathbb{R}/\mathbb{Z}) \), we now use the fact that a \( \text{pin}^- \) structure on a surface is given by a \( \mathbb{Z}/4\mathbb{Z} \)-valued quadratic refinement \( Q : H^1(\Sigma, \mathbb{Z}/2\mathbb{Z}) \to \mathbb{Z}/4\mathbb{Z} \) satisfying
\[
Q(a + b) - Q(a) - Q(b) = 2 \int_\Sigma a \cup b \tag{3.17}
\]
and \( Q(0) = 0 \), see [KT90] for this and other basic properties of \( \text{pin}^- \) structures. Then the evaluation is given by
\[
e^{2\pi \sqrt{-1} \phi([f: \Sigma \to BG])} := \sqrt{-1}^{Q(f^*(\varphi))} e^{2\pi \sqrt{-1} f^*(\alpha)}. \tag{3.18}
\]

Finite-group gauge theories on \( \text{pin}^- \) surfaces for the case when the twisted group superalgebra is the Clifford algebra \( \text{Cl}(n) \) we saw in Lemma 2.7 were studied e.g. in [DG18, Tur18, Kob19].
Again the discussion there can be readily generalized, given our super Frobenius-Schur indicator. The final result when $\Sigma$ is non-orientable is

\[
\frac{1}{|G|} \sum_{f: \pi_1(\Sigma) \to G} \sqrt{-1}^{Q(f^*(\varphi))} e^{2\pi i f\Sigma f^*(\alpha)} = \\
\sum_{\rho \in \text{Irr}(\mathbb{C}[G], \phi, \alpha)} S(\rho)^{\text{ABK}(\Sigma)} \left( \frac{|G|}{\text{qdim} \rho} \right)^{-e(\Sigma)},
\]

where the Arf-Brown-Kervaire invariant $\text{ABK}(\Sigma) \in \{0, 1, \ldots, 7\}$ of the pin$^-$ surface is the class $[\Sigma] \in \Omega^{\text{pin}}_{d=2} = \mathbb{Z}/8\mathbb{Z}$. The defining equation of the super Frobenius-Schur indicator in the statement of Theorem 1.2 is the literal path integral expression of the partition function of the gauge theory on the Möbius strip, when the state on the boundary is specified by $\rho$.

REFERENCES

[Ati71] M. F. Atiyah, *Riemann surfaces and spin structures*, Annales Scientifiques de L’École Normale Supérieure 4 (1971) 47–62.

[BG04] D. Bump and D. Ginzburg, *Generalized Frobenius-Schur numbers*, J. Algebra 278 (2004) 294–313.

[BM16] G. Brumfiel and J. Morgan, *The Pontrjagin Dual of 3-Dimensional Spin Bordism*, arXiv:1612.02860 [math.GT].

[BM18a] ———, *The Pontrjagin Dual of 4-Dimensional Spin Bordism*, arXiv:1803.08147 [math.GT].

[BM18b] ———, *Quadratic Functions of Cocycles and Pin Structures*, arXiv:1808.10484 [math.AT].

[BT13] J. W. Barrett and S. O. G. Tavares, *Two-Dimensional State Sum Models and Spin Structures*, Commun. Math. Phys. 336 (2015) 63–100, arXiv:1312.7561 [math.QA].

[BWHV17] N. Bultinck, D. J. Williamson, J. Haegeman, and F. Verstraete, *Fermionic projected entangled-pair states and topological phases*, Journal of Physics A: Mathematical and Theoretical 51 (2017) 025202, arXiv:1707.00470 [cond-mat.str-el].

[Del99] P. Deligne, *Notes on spinors*, Quantum fields and strings: a course for mathematicians, Vol. 1, 2 (Princeton, NJ, 1996/1997), Amer. Math. Soc., Providence, RI, 1999, pp. 99–135. http://publications.ias.edu/sites/default/files/79_NotesOnSpinors.pdf.

[DG18] A. Debray and S. Gunningham, *The Arf-Brown TQFT of pin$^-$ surfaces*, Topology and quantum theory in interaction, Contemp. Math., vol. 718, Amer. Math. Soc., Providence, RI, 2018, pp. 49–87. arXiv:1803.11183 [math-ph].

[DW90] R. Dijkgraaf and E. Witten, *Topological Gauge Theories and Group Cohomology*, Commun. Math. Phys. 129 (1990) 393.

[FM04] D. S. Freed and G. W. Moore, *Setting the Quantum Integrand of M-theory*, Commun. Math. Phys. 263 (2006) 89–132, arXiv:hep-th/0409135.

[FQ91] D. S. Freed and F. Quinn, *Chern-Simons Theory with Finite Gauge Group*, Commun. Math. Phys. 156 (1993) 435–472, arXiv:hep-th/9111004.

[Fre19] D. S. Freed, *Lectures on Field theory and Topology*, CBMS Regional Conference Series in Mathematics, vol. 133, American Mathematical Society, Providence, RI, 2019. https://bookstore.ams.org/cbms-133/.

[FS06] G. Frobenius and I. Schur, *Über die reellen Darstellungen der endlichen Gruppen*, Sitzungsberichte der Königlich Preußischen Akademie der Wissenschaften zu Berlin (1906) 186–208.

[FS01] J. Fuchs and C. Schweigert, *Category theory for conformal boundary conditions*, Vertex operator algebras in mathematics and physics (Toronto, ON, 2000), Fields Inst. Commun., vol. 39, Amer. Math. Soc., Providence, RI, 2003, pp. 25–70. arXiv:math.CT/0106050.
[GI18] P. Georgieva and E.-N. Ionel, A klein tqft: the local real gromov-witten theory of curves, arXiv:1812.02505 [math.SG].

[GK15] D. Gaiotto and A. Kapustin, Spin TQFTs and Fermionic Phases of Matter, Int. J. Mod. Phys. A31 (2016) 1645044, arXiv:1505.05856 [cond-mat.str-el].

[GOP+18] M. Guo, K. Ohmori, P. Putrov, Z. Wan, and J. Wang, Fermionic Finite-Group Gauge Theories and Interacting Symmetric/Crystalline Orders via Cobordisms, Commun. Math. Phys. (2020) 1–82, arXiv:1812.11959 [hep-th].

[Gow79] R. Gow, Real-valued and 2-rational group characters, J. Algebra 61 (1979) 388–413.

[Gun12] S. Gunningham, Spin Hurwitz numbers and topological quantum field theory, Geom. Topol. 20 (2016) 1859–1907, arXiv:1201.1273 [math.QA].

[Joh80] D. Johnson, Spin structures and quadratic forms on surfaces, Journal of the London Mathematical Society 2 (1980) 365–373.

[KM90] N. Kawanaka and H. Matsuyama, A twisted version of the Frobenius-Schur indicator and multiplicity-free permutation representations, Hokkaido Math. J. 19 (1990) 495–508.

[Kob19] R. Kobayashi, Pin TQFT and Grassmann Integral, JHEP 12 (2019) 014, arXiv:1905.05902 [cond-mat.str-el].

[KT90] R. C. Kirby and L. R. Taylor, Pin structures on low-dimensional manifolds, Geometry of Low-Dimensional Manifolds, vol. 2, London Mathematical Society Lecture Note Series, vol. 151, 1990, pp. 177–242.

[KTTW14] A. Kapustin, R. Thorngren, A. Turzillo, and Z. Wang, Fermionic Symmetry Protected Topological Phases and Cobordisms, JHEP 12 (2015) 052, arXiv:1406.7329 [cond-mat.str-el].

[Med78] A. D. Mednyh, Determination of the number of nonequivalent coverings over a compact Riemann surface, Dokl. Akad. Nauk SSSR 239 (1978) 269–271. English translation: Soviet Math. Doklady 19 (1978), 318–320.

[NR14] S. Novak and I. Runkel, State Sum Construction of Two-Dimensional Topological Quantum Field Theories on Spin Surfaces, J. Knot Theor. Ramifications 24 (2015) 1550028, arXiv:1402.2839 [math.QA].

[Ser16] J.-P. Serre, Finite groups: an introduction, Surveys of Modern Mathematics, vol. 10, International Press, Somerville, MA; Higher Education Press, Beijing, 2016. With assistance in translation provided by Garving K. Luli and Pin Yu.

[Sny07] N. Snyder, Mednykh’s formula via lattice topological quantum field theories, Proc. Centre Math. Appl. Austral. Nat. Univ. 46 (2017) 389–398, arXiv:math.QA/0703073. Proceedings of the 2014 Maui and 2015 Qinhuangdao conferences in honour of Vaughan F. R. Jones’ 60th birthday.

[Tur92] A. Turull, The Schur index of projective characters of symmetric and alternating groups, Ann. of Math. (2) 135 (1992) 91–124.

[Tur07] V. Turaev, Dijkgraaf-Witten invariants of surfaces and projective representations of groups, J. Geom. Phys. 57 (2007) 2419–2430, arXiv:0706.0160 [math.GT].

[Tur18] A. Turzillo, Diagrammatic State Sums for 2D Pin-Minus TQFTs, arXiv:1811.12654 [math.QA].

[Wal64] C. T. C. Wall, Graded Brauer groups, J. Reine Angew. Math. 213 (1963/64) 187–199.

[Wit91] E. Witten, On Quantum Gauge Theories in Two-Dimensions, Commun.Math.Phys. 141 (1991) 153–209.

[WOP+18] J. Wang, K. Ohmori, P. Putrov, Y. Zheng, Z. Wan, M. Guo, H. Lin, P. Gao, and S.-T. Yau, Tunneling Topological Vacua via Extended Operators: (Spin-)TQFT Spectra and Boundary Deconfinement in Various Dimensions, PTEP 2018 (2018) 053A01, arXiv:1801.05416 [cond-mat.str-el].

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