Noncommutative Generalization of Wilson Lines

August 19, 2014

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Abstract

A classical Wilson line is a correspondence between closed paths and elements of a gauge group. However the noncommutative geometry does not have closed paths. But noncommutative geometry have good generalizations of both: the covering projection, and the group of covering transformations. These notions are used for a construction of noncommutative Wilson lines. Wilson lines can also be constructed as global pure gauge fields on the universal covering space. The noncommutative analog of this construction is also developed.

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Foreword

At 1970s I had a strong wish for understanding of physics. I had been reading many books written by physicists, but I did not understand them. Later I started to study math and then I became to understand physics. General relativity became very clear after knowledge of Riemann geometry. Physicist uses the "vector" notion, but mathematician uses rigorous notion of a vector space. It is clear what space mathematicians mean: Banach, Hilbert or Fréchet. A gauge field in physics is denoted by $A_i$, math notation is $\nabla : E \to E \otimes \Omega^1 A$ and sense of math notation is strictly definite. I think that written by physicists texts are accessible to understand for people which have contacts with physicists. I had no contacts with physicists, and so I did not understand their texts. Moreover notation of noncommutative geometry is much more clear then notation of commutative one. For example Levi-Civita connection is given by Christoffel symbols

$$\Gamma^l_{jk} = \frac{1}{2} \sum_r g^{lr} \left\{ \partial_k g_{ri} + \partial_j g_{rk} - \partial_l g_{jk} \right\}.$$ 

But noncommutative notation is much more clear. Written by physicists texts contains a lot of heuristic formulas and it is difficult to distinguish them from rigorous ones. I did not find articles devoted to noncommutative Wilson lines which use math notation. So I have used both physical notation and math one. The correspondence between both these notations is explained. This article has rather math style, so I assigned the mathematics category to it.

1 Introduction

In the commutative gauge theory the curvature of a connection is not sufficient to extract the complete gauge invariant information of the connection. This is the well-known gauge copy problem. For certain gauge groups, the gauge copy problem can be solved by considering the set of all Wilson loops as the basic observables of the gauge theory, see e.g. [5] [11]. There are examples showing that there is also a gauge copy problem in noncommutative geometry, see e.g. Proposition 4.2 from [14]. This means that considering only observables which are constructed from the curvature, we can in general not extract the complete gauge invariant information of the connection. This calls for a suitable generalization of parallel transport and Wilson loops to noncommutative geometry. Some
examples of noncommutative Wilson lines are described in [1, 2, 8]. However these articles contain rather particular examples than a general theory. These examples can be used for the noncommutative torus only. General theory requires a noncommutative generalization of closed paths. However closed paths can be replaced with covering projections. A generalization of covering projections is described in my articles [6, 7]. Following table contains necessary ingredients and their noncommutative analogues.

| Differential geometry | Noncommutative geometry |
|-----------------------|-------------------------|
| Spin manifold [13]    | Spectral triple [16]    |
| Vector bundle [9]     | Projective module [16]  |
| Gauge field [12]      | Noncommutative connection [4, 13] |
| Parallel transport [9] | Noncommutative parallel transport [4, 14] |
| Closed path [15]      | Noncommutative covering projection [6, 7] |

Composition of these ingredients supplies a noncommutative generalization of Wilson lines.

Following notation is used in this article.

| Symbol | Meaning |
|--------|---------|
| Aut(A) | Group * - automorphisms of C*-algebra A |
| C (resp. R) | Field of complex (resp. real) numbers |
| C(X) | C* - algebra of continuous complex valued functions on compact topological space X |
| C_b(X) | C* - algebra of bounded continuous complex valued functions on locally compact topological space X |
| C_0(X) | C* - algebra of continuous complex valued functions on locally compact topological space which tends to 0 at infinity X |
| M(A) | Multiplier algebra of C*- algebra A |
| N | The set of natural numbers |
| Q | The field of rational numbers |
| U(A) ∈ A | Group of unitary operators of algebra A |
| Z | Ring of integers |
| Z_m | Ring of integers modulo m |

2 Prototype. Commutative Wilson Lines

2.1. **Standard description of Wilson lines** [12]. Here the physical notation is used. Let M be a compact Riemann manifold, and let A be a locally pure gauge field, i.e. A can be locally represented by following way

\[ A_i = \partial_i U \cdot U^{-1} \]  \hspace{1cm} (1)

where U is the gauge transformation. Whether (1) is also true globally is another story. Condition (1) is equivalent to that the field stength equals to zero [12]. If the manifold is
not simply connected, π₁(M) ≠ {e} then there is a more general possibility that appears in electrodynamics as the Bohm-Aharonov effect. Let γ be a noncontractible loop in M, beginning and ending at the same point x. Then "Wilson line"

\[ U_γ = P \exp \oint_{γ} A \cdot dx \]  (2)

is gauge invariant, and if \( U_γ \neq 1 \) it cannot be set to one by gauge transformation, \( U_γ \) depends only on \( [γ] \in π_1(M) \). If G is the gauge transformation group then is a group homomorphism

\[ φ : π_1(M) → G, \]  (3)

\[ [γ] ↦ U_γ. \]

2.2. Alternative description of Wilson lines. In this construction I follow to [12] 16.4.1. Let \( \tilde{M} \) be a simply connected manifold, and let F be a discrete symmetry group which acts freely. Let \( M = \tilde{M}/F \). An ordinary field \( ψ(x) \) is equivalent to a field on \( \tilde{M} \) that obeys

\[ ψ(fx) = ψ(x); \ ∀f ∈ F. \]  (4)

Then we can generalize (4) as follows. There is a natural isomorphism \( π_1(M) ≈ F \), and from (3) it follows that there is a natural homomorphism

\[ F → G, \]  (5)

\[ f ↦ U_f. \]

Now we require that \( ψ \) obey not (4) but

\[ ψ(fx) = U_fψ(x); \ ∀f ∈ F. \]  (6)

This operation enables us replace gauge field (2) by "twist" in boundary conditions (6). The gauge field \( A \) which obeys (4) (resp. field \( ψ \)) is replaced with the trivial gauge field (resp. field \( ψ' \) which obey (2)). Field \( ψ' \) is given by

\[ ψ'(y) = \left( P \exp \int_{ω} A' \cdot dx \right) ψ(π(y)) \]

where

- \( π : \tilde{M} → M \) is a covering projection.
- \( ω : [0, 1] → \tilde{M} \) is such that \( ω(0) = y_0 \) is a fixed point, \( ω(1) = y \).
- \( A' \) is a lift of \( A \) by \( π \).

This construction is similar to passive/active approach to physical transformations. Change of a gauge field is similar to a passive transformation of a coordinate system, a change of field is similar to an active transformation of point’s position. Both transformations describe the same physical phenomenon.
2.3. Replacement of closed paths by covering projections. Construction from [12] can be generalized such that the universal covering is replaced by a covering $\pi : \tilde{M} \to M$ such that the image of the composition

$$\pi_1(\tilde{M}) \xrightarrow{\pi_1(\tilde{\pi})} \pi_1(M) \xrightarrow{\varphi} G$$

coincides with image of $\varphi$. In this case homomorphism (3) can be replaced with

$$G(\tilde{M}|M) \to G$$

where $G(\tilde{M}|M)$ is a group of covering transformations [15].

3 Noncommutative Parallel Transport

3.1 Parallel transports

3.1. Definitions of section 2 cannot be directly used in the noncommutative case because the noncommutative geometry does contain closed paths. However paths can be replaced with module parallel transports.

Definition 3.2. [14] Let $A$ be an associative and unital algebra and $E$ a right $A$-module.

1.) A one-parameter group of automorphisms of $A$ is a map $\varphi : \mathbb{R} \times A \to A$, $(\tau, a) \mapsto \varphi(\tau, a) = \varphi_\tau(a)$, such that

(i) $\varphi_\tau(ab) = \varphi_\tau(a)\varphi_\tau(b)$, for all $\tau \in \mathbb{R}$ and $a, b \in A$

(ii) $\varphi_0 = \text{id}_A$

(iii) $\varphi_{\tau+\sigma} = \varphi_\tau \circ \varphi_\sigma$, for all $\tau, \sigma \in \mathbb{R}$

2.) Let $\varphi : \mathbb{R} \times A \to A$ be a one-parameter group of automorphisms of $A$. A module parallel transport on $E$ along $\varphi$ is a map $\Phi : \mathbb{R} \times E \to E$, $(\tau, s) \mapsto \Phi(\tau, s) = \Phi_\tau(s)$, such that

(i) $\Phi_\tau(sa) = \Phi_\tau(s)\varphi_\tau(a)$, for all $\tau \in \mathbb{R}, s \in E$ and $a \in A$

(ii) $\Phi_0 = \text{id}_E$

(iii) $\Phi_{\tau+\sigma} = \Phi_\tau \circ \Phi_\sigma$, for all $\tau, \sigma \in \mathbb{R}$.

If $A$ and $E$ are equipped with a smooth structure, the maps $\varphi$ and $\Phi$ are required to be smooth. Denote by $\text{Trans}_E$ the set of module parallel transports.

3.3. There are different notions of connections [4, 10, 14], some of them are compared in [10]. The space of all connections on $E$ is denoted by $\text{Con}_A(E)$.

Definition 3.4. Let $A$ be a $C$-algebra, and let $E$ be a finite projective $A$ module. Suppose that $\text{Paths}_A$ is a set of one-parameter group of $A$ automorphisms. A connection transport procedure is a natural map

$$\text{Transport} : \text{Con}_A(E) \times \text{Paths}_A \to \text{Trans}_E$$

such that $\text{Transport}(\nabla, \varphi)$ is a transport along $\varphi$ for any $\nabla \in \text{Con}_A(E)$ and $\varphi \in \text{Paths}_A$. 

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Remark 3.5. A connection transport procedure should have a good math and/or physical sense. Such procedures are known in following cases:

1. Commutative differential geometry [9, 14].
2. \( A = \mathbb{M}_n(\mathbb{C}) \) [14].
3. Noncommutative torus [1, 2, 5, 11].
4. Noncommutative differential geometry [4].

3.2 Parallel transform in noncommutative differential geometry

In this section I follow to [4].

3.6. [4] Let \((A, G, \alpha)\) be a \( C^* \) dynamical system, where \( G \) is a Lie group. We shall say that \( x \in A \) is of \( C^\infty \) class iff the map \( g \mapsto \alpha_g(x) \) from \( G \) to the normed space \( A \) is in \( C^\infty \). The involutive algebra \( A^\infty = \{ x \in A, x \text{ of class } C^\infty \} \) is norm dense in \( A \).

Let \( E^\infty \) be a finite projective module on \( A^\infty \), (we shall write it as a right module); \( E = E^\infty \otimes_{A^\infty} A \) is then a finite projective module on \( A \).

Lemma 3.7. [4] For every finite projective module \( E \) on \( A \), there exists a finite projective module \( E^\infty \) on \( A^\infty \), unique up to isomorphism, such that \( E \) is isomorphic to \( E^\infty \otimes_{A^\infty} A \).

3.8. Let \( \delta \) be the representation of \( \text{Lie} \, G \) in the Lie-algebra of derivations of \( A^\infty \) given by

\[
\delta_X(x) = \lim_{t \to 0} \frac{1}{t} (\alpha_{\delta_t}(x) - x), \quad \text{where} \quad \delta_0 = X, \, x \in A^\infty.
\]

Definition 3.9. [4] \( E^\infty \) be a finite projective module on \( A^\infty \), a connection (on \( E^\infty \)) is a linear map \( \nabla : E^\infty \to E^\infty \otimes (\text{Lie} \, G)^* \) such that, for all \( X \in \text{Lie} \, G \) and \( \xi \in E^\infty, \, x \in A^\infty \) one has

\[
\nabla_X(\xi \cdot x) = \nabla_X(\xi) \cdot x + \xi \cdot \delta_X(x).
\]

3.10. There is a natural correspondence between elements of \( \text{Lie} \, G \) and one-parameter transformation groups [9]. Let \( \nabla \) be a connection (on \( E^\infty \)). If \( X \in \text{Lie} \, G \) defines a one-parameter group of automorphisms \( \varphi \) of \( A^\infty \) then \( \nabla_X \) defines a module parallel transport \( \Phi \) on \( E^\infty \) along \( \varphi \). So we have a connection transport procedure

\[
\text{Transport} : \text{Con}_{A^\infty}(E^\infty) \times \text{Paths}_{A^\infty} \to \text{Trans}_{E^\infty}.
\]

Definition 3.11. [4] Let \( \nabla \) be a connection on the finite projective module \( E^\infty \) on \( A^\infty \), the curvature of \( \nabla \) is the element \( \mathcal{T} \) of \( \text{End}_{A^\infty}(E^\infty) \otimes \mathbb{L}^2(\text{Lie} \, G)^* \) given by

\[
\mathcal{T}(X, Y) = \nabla_X \nabla_Y - \nabla_Y \nabla_X - \nabla_{[X,Y]} \in \text{End}_{A^\infty}(E^\infty), \quad \forall X, Y \in \text{Lie} \, G.
\]

Definition 3.12. [4] A connection with zero curvature is said to be flat. Denote by \( \text{Con}_{A}(E)_0 \) a space of flat connections.

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If \( e \in A^\infty \) is an idempotent then every connection on \( e A^\infty \) is of the form \( \nabla_X(e) = \nabla_X^0(e) + \theta_X e \), \( \forall X \in e A^\infty \), \( X \in \text{Lie} \, G \), where the form \( \theta \in e \Omega^1 e \) is uniquely determined by \( \nabla \). One has \( \theta_X = -\theta_X \), \( \forall X \in \text{Lie} \, G \) iff \( \nabla \) is compatible with the hermitian structure of \( e A^\infty \). We identify \( \text{End}(e A^\infty) \) with \( e A^\infty \subset A^\infty \), the curvature \( T_0 \) of the grassmannian connection is the 2-form \( e(d\theta \wedge \theta) \in \Omega^2 \), the curvature of \( \nabla = \nabla^0 + \theta \wedge \) equals to

\[
T_0 + e(d\theta + \theta \wedge \theta) e \in \Omega^2.
\]

## 4 Noncommutative Generalization of Closed Paths

### 4.1 Commutative loops

Some of parallel transports can be regarded as noncommutative loops. First of all we consider the commutative case. Let \( M \) be a manifold and let \( \pi : \tilde{M} \to M \) be a covering projection, \( G(M|M) \) is the group of covering transformations, i.e. \( M \approx \tilde{M}/G(M|M) \). Let \( \varphi : \mathbb{R} \times C^\infty(M) \to C^\infty(M) \) satisfies condition 1 of definition 3.2.

From Gelfand - Naimark theorem \([3]\) it follows that \( \varphi \) defines a one-parameter group of homeomorphisms (indeed diffeomorphisms) \( \varphi^* : \mathbb{R} \otimes M \to M \). If \( \varphi_1 = \text{Id}_{C^\infty(M)} \) then a function \( \gamma : [0,1] \to M \), \( t \mapsto \varphi^*(t,x_0) \) corresponds to a closed path. We would like to know whether this path is not contractible. It is known \([15]\) that covering projections have property of unique path lifting. So \( \varphi^* \) can be lifted to \( \tilde{\varphi}^* : \mathbb{R} \times C^\infty(\tilde{M}) \to C^\infty(\tilde{M}) \). If \( \tilde{\varphi}_1^* \in G(\tilde{M}|M) \) then \( \tilde{\varphi}_1^* = \text{Id}_M \). If \( \tilde{\varphi}_1^* \) is not a trivial element in \( G(\tilde{M}|M) \) then \( \varphi \) does not correspond to a contractible path.

### 4.2 Noncommutative loops

In my articles \([6,7]\) there is a construction of covering projections for \( C^\ast \)-algebras (See Appendix 8.1). A covering projection of \( C^\ast \)-algebra \( A \) is a \( * \)-homomorphism \( A \to M(\tilde{A}) \) where \( \tilde{A} \) is another \( C^\ast \)-algebra. A group of noncommutative covering transformations \( G(\tilde{A}|A) \) acts on \( \tilde{A} \) such that \( g(a\tilde{a}) = ag\tilde{a} \), \( \forall a \in A \), \( \forall a \in \tilde{A} \) \( \forall g \in G(\tilde{A}|A) \). For any \( \varphi : \mathbb{R} \times A \to A \) there is an unique lift \( \tilde{\varphi} : \mathbb{R} \times \tilde{A} \to \tilde{A} \) such that

\[
\tilde{\varphi}_\tau(a\tilde{a}) = \varphi(\tau)(a)\tilde{\varphi}_\tau(a); \quad a \in A, \quad \tilde{a} \in \tilde{A}, \quad \tau \in \mathbb{R}.
\]

**Definition 4.3.** If \( \varphi \) is such that

\[
\tilde{\varphi}_1 = g \in G(\tilde{A}|A),
\]

\[
\tilde{\varphi}_\tau \notin G(\tilde{A}|A), \quad 0 < \tau < 1.
\]

then we say that \( \varphi \) is a closed path associated with \( g \).

**Definition 4.4.** Suppose that for any \( g \in G(\tilde{A}|A) \), \( \nabla \in \text{Con}_A(\mathcal{E})_0 \) and for any closed paths \( \varphi', \varphi'' \) associated with \( g \) we have

\[
\text{Transport}(\nabla, \varphi'_1) = \text{Transport}(\nabla, \varphi''_1).
\]

A generalized Wilson line is a map

\[
\text{Wilson} : G(\tilde{A}|A) \times \text{Con}_A(\mathcal{E})_0 \to \text{Aut}(\mathcal{E}),
\]

\[
(g, \nabla) \mapsto \text{Transport}(\nabla, \varphi_1)
\]

where \( \varphi \) is associated with \( g \) and \( \text{Transport} \) is defined in \([5.10]\)
5 Wilson Lines and Noncommutative Covering Spaces

Commutative geometry has a lot of local structures, for example local sections of bundles. There are bundles such that they have no global sections. Let $p : P \to M$ be a bundle such that $p$ has no global sections. Let $\pi : \tilde{M} \to M$ be the universal covering projection, and $\tilde{p} : \tilde{M} \to M$ be the pullback \([15]\) of $p$ by $\pi$. Then $\tilde{p}$ can have global sections. Any locally pure gauge field can be locally represented by (1), but cannot be represented by (1) globally in general case. However pullback of this field can be globally represented by (1). Noncommutative geometry has no local sections. However there is the noncommutative generalisation of covering projections \([6, 7]\). Local pure gauge fields can be regarded as global pure gauge fields on universal noncommutative covering projections. Locally gauge fields satisfy (1) which can be rewritten in physical notation \([1]\)

$$\partial_i g = iA_i * g. \quad (9)$$

This equation has a noncommutative analog. Let $A^\infty$ be a smooth algebra and $\mathcal{E}^\infty$ be a finitely generated projective smooth $A^\infty$ module. Let $(A^\infty, \tilde{A}^\infty, \tilde{G}, \tilde{A}^\infty_M)$ be a noncommutative covering projection (See Appendix \([8, 1]\). Spaces $A^\infty, \tilde{A}^\infty, \mathcal{E}^\infty$ are operator spaces. Let denote $\tilde{\mathcal{E}}^\infty = \mathcal{E}^\infty \otimes_{A^\infty} \tilde{A}^\infty$ where $\otimes$ means the Haagerup tensor product. The $\tilde{\mathcal{E}}^\infty$ module can be regarded as pullback of $\mathcal{E}^\infty$. Any vector $X \in \text{Lie} \tilde{G}$ can be lifted to the vector $\tilde{X} \in \text{Lie} \tilde{G}$ where $\text{Lie} \tilde{G}$ is a Lie algebra of infinitesimal transformations of the $\tilde{A}^\infty$. Any connection $\nabla : \mathcal{E}^\infty \to \mathcal{E}^\infty \otimes \Omega^1 A^\infty$ can be lifted to $G$ equivariant connection $\tilde{\nabla} : \tilde{\mathcal{E}}^\infty \to \tilde{\mathcal{E}}^\infty \otimes \Omega^1 \tilde{A}^\infty$. Otherwise any connection $\tilde{\nabla} : \tilde{\mathcal{E}}^\infty \to \tilde{\mathcal{E}}^\infty \otimes \Omega^1 \tilde{A}^\infty$ naturally induces a map $\tilde{\nabla}^l : \text{End}_{\tilde{A}^\infty}(\tilde{\mathcal{E}}^\infty) \to \text{End}_{\tilde{A}^\infty}(\tilde{\mathcal{E}}^\infty) \otimes \Omega^1 \tilde{A}^\infty$. Noncommutative analog of (9) is given by

$$\tilde{X}U = \tilde{\nabla}_X U; \forall X \in \text{Lie} \tilde{G}, U \in \text{Aut}(\tilde{\mathcal{E}}^\infty). \quad (10)$$

$U$ is a noncommutative analog of a global gauge transformation on the universal covering space. Wilson line can be regarded as a group homomorphisms $G \to \text{Aut}(\tilde{\mathcal{E}}^\infty)$ given by

$$g \mapsto (gU) \cdot U^{-1} \quad (11)$$

where $g \in G(\tilde{A}^\infty | A^\infty)$ and $U$ satisfies (10).

6 Alternative Description of Wilson Lines

As well as in \([2, 2]\) (See \([12]\)) we can define an alternative approach to Wilson lines. Suppose that there is a spectral triple $(\mathcal{A}, H, D)$ \([16]$. In \([6, 7] I$ defined noncomutative covering projection $(\mathcal{A}, H, D)$ with a group of covering transformations $G(\mathcal{A}, \mathcal{A})$, such that there is the natural *-homomorphism $\mathcal{A} \to M(\tilde{A})$ and representation $\pi : \mathcal{A} \to U(\tilde{H})$. Suppose that $\rho : G(\mathcal{A}, \mathcal{A}) \to U(\tilde{H})$ is a representation such that

$$\tilde{D}\rho(g)h = \rho(g)\tilde{D}h, \forall h \in \text{Dom}(\tilde{D}).$$

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Let $\tilde{H}_0 = \{ h \in \tilde{H} \mid \pi(g)h = \rho(g)h \}$. Algebra $A$ naturally acts on $\tilde{H}_0$ and we have a twisted spectral triple $(\tilde{A}, \tilde{H}_0, \tilde{D})$ which can be regarded as description of Wilson line.

### 7 Wilson Lines on the Noncommutative Torus

#### 7.1. A flat connection

Let $\theta \in [0,1] - Q$ and $A_\theta$ the $C^*$-algebra generated by two unitaries $u, v$ such that $uv = \lambda vu, \lambda = \exp(2\pi i \theta)$ which is said to be a noncommutative torus. There is a pre-$C^*$-algebra $A_\theta$ of smooth functions defined in [16]. There are one-parameter groups $\phi^u, \phi^v$ such that

\begin{align}
\phi^u_t(u) &= e^{2\pi itu}, \quad \phi^u_t(v) = v, \\
\phi^v_t(u) &= u, \quad \phi^v_t(v) = e^{2\pi itv}.
\end{align}

Let $E = A_\theta$. Let $\omega = i(c_\mu du + c_\nu d\nu) \in \Omega^1 A_\theta$ be such that $c_\mu, c_\nu \in \mathbb{R}$ then $\omega^X = -\omega^X, \forall X \in \text{Lie} G$. From [3.13] it follows that there is a connection connection $\nabla$ given by

$$a \mapsto a \otimes \omega; \quad a \in E = A.$$

From (13) it follows that curvature of $\nabla$ is equal to

$$1(d1 \wedge d1) + dw + w \wedge w.$$ 

All summands of above equation equal to zero, so curvature of $\nabla$ is zero, i.e. $\nabla$ is flat. One parameter groups $\phi^u, \phi^v$ correspond to following module parallel transports.

\begin{align}
\Phi^u_t(a) &= e^{ic_\nu \tau} a, \\
\Phi^v_t(a) &= e^{ic_\mu \tau} a.
\end{align}

**Example 7.2.** In my article [6] I have found all noncommutative covering projections of $A_\theta$. Let $\theta' = \theta/4$ and $A_{\theta'}$ is generated by $x, y$ such that $xy = e^{2\pi i \theta'}yx$. There is a $*$-homomorphism $\pi : A_\theta \rightarrow A_{\theta'}$ such that

$$\pi(u) = x^2, \quad \pi(v) = y^2.$$ 

(14)

The group of transformation coverings $G(A_{\theta'} | A_\theta)$ is isomorphic to $\mathbb{Z}_2 \times \mathbb{Z}_2$ with two generators $g_u, g_v \in \mathbb{Z}_2 \times \mathbb{Z}_2$ such that

$$g_u x = -x, \quad g_u y = y, \quad g_v x = x, \quad g_v y = -y.$$ 

(15)

One-paramemeter groups $\phi^u$ and $\phi^v$ can be lifted to $\bar{\phi}^u$ and $\bar{\phi}^v$ such that

\begin{align}
\bar{\phi}^u_t(x) &= e^{2\pi itu} x, \quad \bar{\phi}^u_t(y) = y, \\
\bar{\phi}^v_t(x) &= x, \quad \bar{\phi}^v_t(y) = e^{2\pi itv} y.
\end{align}

From above equations it follows that $\bar{\phi}^u_1 = g_u, \bar{\phi}^v_1 = g_v$. So $\phi^u$ and $\phi^v$ are closed paths associated with $g_u$ and $g_v$ respectively. Let $\nabla$ be a flat connection defined in [7.1] Generalized Wilson line is given by

$$\text{Wilson}(g_u, \nabla)(a) = e^{2\pi i tu} a,$$

$$\text{Wilson}(g_v, \nabla)(a) = e^{2\pi i tv} a.$$ 

where $a \in E$. 

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Let $\tilde{A}_0$ be an infinite covering projection of noncommutative torus (See 5.2). It is known that the covering transformation group equals to $\mathbb{Z}^2$. Let us consider an unitary operator $U \in M\left(\tilde{A}_0\right)$ given by

$$U = \pi_{\mu}(\phi^u) \pi_{\nu}(\phi^v)$$

where functions $\phi^u, \phi^v \in C_b(\mathbb{R})$ are given by

$$\phi^u(x) = e^{ic_u x}, \phi^v(x) = e^{ic_v x}, \forall x \in \mathbb{R}.$$  

It is easy to show that $U$ satisfies to (10). A group homomorphism (11) is given by

$$n_1 \mapsto e^{2\pi i c_1} 1_{M(\tilde{A}_0)}, n_2 \mapsto e^{2\pi i c_2} 1_{M(\tilde{A}_0)}$$

where $n_1, n_2 \in \mathbb{Z}^2$ are generators of the covering transformation group.

7.3. Let $E = A^1$ be a free module and let $e_1, \ldots, e_4 \in E$ be its generators. Let $\nabla : E \to E \otimes \Omega^1(A_0)$ be a connection given by

$$\nabla e_1 = c_u e_2 \otimes du, \nabla e_2 = -c_u e_1 \otimes du, \nabla e_3 = c_v e_4 \otimes dv, \nabla e_4 = -c_v e_3 \otimes dv.$$

where $c_u, c_v \in \mathbb{R}$. Let $X, Y \in \text{Lie}G$ correspond to one-parameters groups given by (12) and (13) respectively. We have $[X, Y] = 0$ because given by (12) and (13) groups commute. A direct calculation shows that $\nabla_X \nabla_Y = \nabla_Y \nabla_X$, so we have

$$\mathcal{T}(X, Y) = \nabla_X \nabla_Y - \nabla_Y \nabla_X - \nabla_{[X, Y]} = 0.$$  

Since $X$ and $Y$ are generators of $\text{Lie}G$ we have $\mathcal{T} \equiv 0$, i.e a connection $\nabla$ is flat.

**Example 7.4.** Let $\pi : A_0 \to A_0'$ be a noncommutative covering projection given by (14). Then generalised Wilson line is given by

$$\text{Wilson}(g_u, \nabla)(e) = \begin{pmatrix}
\cos(2\pi c_u) 1_{M(\tilde{A}_0)} & -\sin(2\pi c_u) 1_{M(\tilde{A}_0)} & 0 & 0 \\
\sin(2\pi c_u) 1_{M(\tilde{A}_0)} & \cos(2\pi c_u) 1_{M(\tilde{A}_0)} & 0 & 0 \\
0 & 0 & 1_{M(\tilde{A}_0)} & 0 \\
0 & 0 & 0 & 1_{M(\tilde{A}_0)}
\end{pmatrix} e,$$

$$\text{Wilson}(g_v, \nabla)(e) = \begin{pmatrix}
1_{M(\tilde{A}_0)} & 0 & 0 & 0 \\
0 & 1_{M(\tilde{A}_0)} & 0 & 0 \\
0 & 0 & \cos(2\pi c_v) 1_{M(\tilde{A}_0)} & -\sin(2\pi c_v) 1_{M(\tilde{A}_0)} \\
0 & 0 & \sin(2\pi c_v) 1_{M(\tilde{A}_0)} & \cos(2\pi c_v) 1_{M(\tilde{A}_0)}
\end{pmatrix} e$$

where $e \in E$. Let $c_u, s_u, c_v, s_v \in C_b(\mathbb{R})$ be such that

$$c_u(x) = \cos(c_u x), s_u(x) = \sin(c_u x), c_v(x) = \cos(c_v x), s_v(x) = \sin(c_v x); \forall x \in \mathbb{R}.$$  

Let \( U \in \text{Aut}_{\tilde{\mathcal{A}}_\theta} (\tilde{\mathcal{E}}) \) be an unitary element given by

\[
U = \begin{pmatrix}
\pi_u (c_u) & -\pi_u (s_u) & 0 & 0 \\
\pi_u (s_u) & \pi_u (c_u) & 0 & 0 \\
0 & 0 & \pi_v (c_v) & -\pi_v (s_v) \\
0 & 0 & \pi_v (s_v) & \pi_v (c_v)
\end{pmatrix}
\]

where \( \pi_u, \pi_v \) are defined in 8.2.
Element \( U \) satisfies to (10). A group homomorphism \( \theta \) is given by

\[
\begin{align*}
n_1 & \mapsto \begin{pmatrix}
\cos(2\pi c_u)1_{M(\tilde{\mathcal{A}}_\theta)} & -\sin(2\pi c_u)1_{M(\tilde{\mathcal{A}}_\theta)} & 0 & 0 \\
\sin(2\pi c_u)1_{M(\tilde{\mathcal{A}}_\theta)} & \cos(2\pi c_u)1_{M(\tilde{\mathcal{A}}_\theta)} & 0 & 0 \\
0 & 0 & 1_{M(\tilde{\mathcal{A}}_\theta)} & 0 \\
0 & 0 & 0 & 1_{M(\tilde{\mathcal{A}}_\theta)}
\end{pmatrix}, \\
n_2 & \mapsto \begin{pmatrix}
1_{M(\tilde{\mathcal{A}}_\theta)} & 0 & 0 & 0 \\
0 & 1_{M(\tilde{\mathcal{A}}_\theta)} & 0 & 0 \\
0 & 0 & \cos(2\pi c_v)1_{M(\tilde{\mathcal{A}}_\theta)} & -\sin(2\pi c_v)1_{M(\tilde{\mathcal{A}}_\theta)} \\
0 & 0 & \sin(2\pi c_v)1_{M(\tilde{\mathcal{A}}_\theta)} & \cos(2\pi c_v)1_{M(\tilde{\mathcal{A}}_\theta)}
\end{pmatrix}.
\end{align*}
\]

\section{Appendix. Noncommutative covering projections}

\subsection{General Theory}

Let us remind notion of noncommutative covering projection.

\textbf{Definition 8.1.} [6]
Let \( \mathcal{A}_X \mathcal{B} \) be a Hermitian \( \mathcal{B} \)-rigged \( \mathcal{A} \)-module, \( G \) is finite or countable group such that

- \( G \) acts on \( \mathcal{A} \) and \( \mathcal{X} \),
- Action of \( G \) is equivariant, i.e \( g(a\xi) = (ga)(g\xi) \), and \( \mathcal{B} \) invariant, i.e \( g(\xi b) = (g\xi)b \) for any \( \xi \in \mathcal{X}, b \in \mathcal{B}, a \in \mathcal{A}, g \in G \),
- Inner-product of \( G \) is equivariant, i.e \( \langle g\xi, g\zeta \rangle_\mathcal{X} = \langle \xi, \zeta \rangle_\mathcal{X} \) for any \( \xi, \zeta \in \mathcal{X}, g \in G \).

Then we say that \( \mathcal{A}_X \mathcal{B} \) is a \( G \)-equivariant \( \mathcal{B} \)-rigged \( \mathcal{A} \)-module.

If \( \mathcal{B}, \mathcal{A}, G, \mathcal{A}_X \mathcal{B} \) satisfy definition 5.6 [7] than we say that quadruple \( (\mathcal{B}, \mathcal{A}, G, \mathcal{A}_X \mathcal{B}) \) is an \textit{infinite noncommutative covering projection}. Finite covering projections are particular cases of infinite ones. If \( (\mathcal{B}, \mathcal{A}, G, \mathcal{A}_X \mathcal{B}) \) is an infinite covering projection than \( G \) acts on \( \mathcal{A} \) there is a \( * \)-homomorphism \( \pi : \mathcal{B} \rightarrow M(\mathcal{A}) \) such that

\[
g(\pi(b)a) = \pi(b)(ga), \ b \in \mathcal{B}, a \in \mathcal{A}, g \in G.
\]  

(16)

If \( \mathcal{B} \in \mathcal{B} \) is a smooth subalgebra then there is a smooth version \( (\mathcal{B}, \mathcal{A}, G, \mathcal{A}_X \mathcal{B}) \) of covering projection defined in section 8 of [7].
8.2 Infinite Noncommutative Covering Projection of Noncommutative Torus

Let $A_\theta$ be a noncommutative torus generated by unitary elements $u, v$. In 7.2 [7] it is constructed an infinite covering projection $(A_\theta, \tilde{A}_\theta, \mathbb{Z}^2, X_{A_\theta})$. This construction contains a representation $A_\theta \to B(H)$ and two representations $\pi_u : C_0(\mathbb{R}) \to B(H)$, $\pi_v : C_0(\mathbb{R}) \to B(H)$ such that $\tilde{A}_\theta$ is the norm completion of generated by following operators

$$\pi_u(f)\pi_v(g), \pi_v(f), \pi_u(g); f, g \in C_0(\mathbb{R}).$$

(17)

This subalgebra of $B(H)$. Both $A_\theta$ and $\tilde{A}_\theta$ are represented in the same Hilbert space $H$. These representations induce a $*$-homomorphism $A_\theta \to M(\tilde{A}_\theta)$ given by

$$u \cdot \pi_u(f)\pi_v(g) = \pi_u(\phi^{\exp f})\pi_v(g),$$

$$v \cdot \pi_u(f)\pi_v(g) = \pi_u(f)\pi_v(\phi^{\exp g}).$$

where $\phi^{\exp} \in C_b(\mathbb{R})$ be given by $x \mapsto e^{ix}$ ($x \in \mathbb{R}$). For any $f \in C_0(\mathbb{R})$ denote by $f^\dagger \in C_0(\mathbb{R})$ given by

$$f^\dagger(x) = f(x + 2\pi); \forall x \in \mathbb{R}.$$  

Noncommutative group of covering transformation equals to $\mathbb{Z}^2$ and generators $n_1, n_2 \in \mathbb{Z}^2$ act on operators (17) by following way

$$n_1 \cdot \pi_u(f)\pi_v(g) = \pi_u(f^\dagger)\pi_v(g), \ n_2 \cdot \pi_u(f)\pi_v(g) = \pi_u(f)\pi_v(g^\dagger);$$

Extension of this action gives an action of $\mathbb{Z}^2$ on $\tilde{A}_\theta$.

References

[1] Anton Alekseev, Andrei Bytsko, Wilson lines on noncommutative tori. arXiv:hep-th/0002101, 2000.

[2] J. Ambjorn, Y. M. Makeenko, J. Nishimura and R. J. Szabo, Lattice gauge fields and discrete noncommutative Yang-Mills theory. JHEP 0005, 023 hep-th/0004147, 2000.

[3] W. Arveson. An Invitation to C*-Algebras, Springer-Verlag. ISBN 0-387-90176-0, 1981.

[4] Alain Connes. C*-algebras and differential geometry. arXiv:hep-th/0101093, 2001.

[5] B. Driver. Classifications of bundle connection pairs by parallel translation and lassos. J. Funct. Anal. 83, no. 1, (1989).

[6] Petr Ivankov. Finite covering projections of noncommutative torus. arXiv:1402.0775, 2014.

[7] Petr Ivankov. Infinite Noncommutative Covering Projections. arXiv:1405.1859, 2014.
[8] N. Ishibashi, S. Iso, H. Kawai and Y. Kitazawa, *Wilson loops in noncommutative Yang-Mills*. Nucl. Phys. B 573, 573 hep-th/9910004, 2000.

[9] Kobayashi S., Nomizu K. *Foundations of Differential Geometry*. Vol. 1.-Wiley-Interscience, 1963.

[10] Partha Sarathi Chakraborty, Satyajit Guin. *Equivalence of Two Approaches to Yang-Mills on Non-commutative Torus*. arXiv:1304.7616v1, 2013.

[11] A. Sengupta. *Gauge invariant functions of connections*. Proc. Am. Math. Soc. 121, 897-905 (1994).

[12] *Superstring Theory: Volume 2, Loop Amplitudes, Anomalies and Phenomenology*. (Cambridge Monographs on Mathematical Physics) by Michael B. Green, John H. Schwarz, Edward Witten. 1988

[13] *Lecture notes on noncommutative geometry and quantum groups*, Edited by Piotr M. Hajac.

[14] Alexander Schenkel. *Module parallel transports in fuzzy gauge theory*, arXiv:1201.4785, 2013.

[15] E.H. Spanier. *Algebraic Topology*. McGraw-Hill. New York 1966.

[16] J.C. Várilly. *An Introduction to Noncommutative Geometry*. EMS 2006.