Role of Hydrogen in the Electronic Properties of CaFeAsH-based Superconductors

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The electronic and magnetic properties of the new hydride superconductor CaFeAsH, which superconducts up to 47 K when electron-doped with La, and the isovalent alloy system CaFeAsH$_{1-x}$F$_x$ are investigated using density functional based methods. The $\bar{Q}= (\pi, \pi, 0)$ peak of the nesting function $\xi(\bar{q})$ is found to be extremely strong and sharp, and additional structure in $\xi(\bar{q})$ associated with the near-circular Fermi surfaces (FSs) that may impact low energy excitations is quantified. The unusual band introduced by H, which shows strong dispersion perpendicular to the FeAs layers, is shown to be connected to a peculiar van Hove singularity just below the Fermi level. This band provides a three dimensional electron ellipsoid Fermi surface not present in other Fe-based superconducting materials nor in CaFeAsF. Electron doping by 25% La or Co has minor affect on this ellipsoid Fermi surface, but suppresses FS nesting strongly, consistent with the viewpoint that eliminating strong nesting and the associated magnetic order allows high T$_c$ superconductivity to emerge. Various aspects of the isovalent alloy system CaFeAsH$_{1-x}$F$_x$ and means of electron doping are discussed in terms of influence of incipient bands.

I. INTRODUCTION

Identifying and studying similar materials that display very different behavior has been an active approach to understanding the origin of high temperature (T$_c$) superconducting Fe-based pnictides and chalcogenides, and is expected to promote the quest for new and possibly higher T$_c$ superconductors. In this class of materials, elemental substitution can be done in an isovalent manner to provide (seemingly) minor changes, while aliovalent substitution supplies either electron or hole carriers to provide (seemingly) minor changes, while aliovalent elemental substitution can be done in an isovalent manner under various conditions, adding a new dimension to the phenomenon. Hosono et al. showed that hydrogen in this material can form a solid solution in any ratio with F, making CaFeAsF$_{1-x}$H$_x$ an isovalent, isostructural system in which specific differences can be varied continuously. Though Muraba et al. have emphasized some important differences from the oxide analogs in the electronic structure, the differences arising from H versus F have not been studied thoroughly. Such isovalent sister compounds provide means to evaluate the effects of specific, seemingly small differences, such as has been exploited in the LiFeAs, NaFeAs, and MgFeGe compounds, which likewise are isostructural and isostructural with extremely similar band structures but differing properties. In fact, MgFeGe is not superconducting at all. Though non-superconducting at ambient pressure, both CaFeAsH and CaFeAsF display pressure-driven superconductivity with T$_c$ up to 25-28 K at the superconducting onset at 3-5 GPa, beyond which T$_c$ decreases with additional pressure.

Impressively, indirect electron-doping by La on the Ca site of CaFeAsH produces superconductivity, with T$_c^{max}$=47 K, in the regime of the rare earth 1111 compounds with the highest T$_c$'s (up to 56 K). This high T$_c$ can be contrasted with “direct” electron doping by Co-substitution of Fe, which attains only T$_c^{max}$=23 K. Also noteworthy is that in both cases there is no superconducting dome – the “onset T$_c$” (some samples are evidently not homogeneous) is almost constant from the appearance of superconductivity at a doping level at $x=0.07$-0.08 to the maximum doping studied, $x=0.3$-0.4.

There are strong parallels to another two dimensional (2D) electron-doped system A$_2$MCl, where $M = $ Zr (respectively Hf), $A$ is an alkali metal, with a T$_c$ of 15 K (respectively 26 K). In these two compounds T$_c$ is independent of doping level except at the low end, at the onset of superconducting where T$_c$ is somewhat higher. This indifference of T$_c$ to the evolving electronic structure has been discussed and this same indifference in CaFeAsH raises important questions, since there is a great deal of evidence in other Fe-based superconductors that properties including T$_c$ are sensitive to small changes in electronic structure. There are differences in T$_c$ and other properties with chemical variation on the H/F site, in addition to the doping differences, that so far have no explanation.

Incorporation of H as an anion in the blocking layer and the realization of the superconductivity as high as 47 K invites certain questions. First, within the same 1111 phase, why does T$_c$ increase so much with La doping after (LaO)$^+$ layers are replaced by isovalent (CaH)$^+$ layers? Second, in comparison with CaFeAsF, what is the role of H in stabilizing such a high T$_c$? A third question is why the maximal superconducting critical temperatures in La-doped and Co-doped CaFeAsH samples (both electron doped) are so different. These are the sort of issues
we try to shed light on.

Using methods similar to those we use (see below), Muraba et al.\cite{1} and Wang and collaborators\cite{10} have identified the main differences in the band structures of CaFeAsH and CaFeAsF: the former has an additional Fermi surface arising from a band that is displaced from its position in the F compound and in other 1111 compounds. While it is in all cases a primarily Fe 3d band, the H 1s states are less strongly bound than the F 2p states and evidently provide a propitious pathway for Fe-As-H-As-Fe hopping from FeAs layer to FeAs layer. Fermi surface nesting was discussed but the difference between the H and F compounds was not quantified.

These differences, and the electronic structure to be analyzed here, call to mind the recent discussion on the possible influence of incipient bands in Fe-based superconductors. Incipient bands are those that do not actually cross the Fermi energy (E_F) and thus do not have an associated Fermi surface, but are near enough to the Fermi level (within a pairing cutoff energy) to contribute to pairing and superconductivity. Chen et al.\cite{11} who provide an excellent overview of the discussion, conclude that incipient bands may indeed play a role in pairing in LiFeAs and in FeSe monolayers on SrTiO_3 substrates. The difference in the electronic structures at or near the Fermi level, and the associated (magnetic) fluctuations, in these H and F compounds must account for their differing properties, and the incipient band viewpoint may be a useful one to consider.

In this paper first principles methods, described in Sec. II, are applied to study the electronic structures and magnetic ground states of CaFeAsH and CaFeAsF, and the non-magnetic and superconducting electron-doped phases Ca_{0.75}La_{0.25}FeAsH and CaFe_{0.75}Co_{0.25}AsH. Sec. III contains the main results. It is confirmed that CaFeAsH has, in addition to the nearly circular, 2D electron and hole cylinders expected of an Fe 1111 compound, a strongly k_z-dispersive band crossing the Fermi level (E_F) that leads to a three dimensional (3D) ellipsoidal Fermi pocket, whose implications we explore. A specific feature that we quantify is extremely sharp \( Q = (\pi, \pi, 0) \) nesting (as usual, in units of in-plane lattice constant a=1). Several other differences between the H- and F-compounds are analyzed, including nesting away from the nesting peak at \( \tilde{Q} \) that may impact low energy properties. A short summary is provided in Sec. IV.

II. STRUCTURE AND METHODS

A. Crystal Structure

The crystal symmetry across the CaFeAsF_{1-x}H_x system is tetragonal P4/nmm, comprised of alternating CaH_2 and Fe_2As_2 layers as in Fe-based 1111 compounds.\cite{12} The hydride and fluorides ions have closed shell, negatively charged ionic configurations with not greatly differing anionic radius, allowing the successful synthesis of all solid solutions.\cite{13} As mentioned, CaFeAsH displays differences compared to CaFeAsF\cite{14,15} which is part of the motivation for this study.

The experimental lattice parameters are used in the calculations. In the P4/nmm phase above the magnetic ordering temperature, for CaFeAs\cite{16} \( a_t = 3.879 \) Å, \( c_t = 8.26 \) Å. For CaFeAs\cite{17} the lattice parameter \( a_t = 3.878 \) Å is nearly identical, while \( c_t = 8.593 \) Å. Thus the larger F ion, confirmed by charge density plots (not shown), enlarges the structure only along the c axis. For CaFeAsH in the magnetically ordered Cmma phase, \( a = 5.457 \) Å = 1.407\( \alpha \), \( b = 5.492 \) Å = 1.416\( \alpha \), and \( c = 8.21 \) Å = 0.994\( \ell \). For Ca_{0.75}La_{0.25}FeAsH and CaFe_{0.75}Co_{0.25}AsH in the tetragonal phase, the unit cell size is 2\( \alpha \) \times 2\( \alpha \) \times c, \( a = 3.883 \) Å = 1.001\( \alpha \), \( c = 8.20 \) Å = 0.993\( \ell \).\cite{17}

It will be useful to be aware of the distances between the As plane and the anion plane. For CaFeAsH, \( d(\text{As-H}) = 2.71 \) Å; for CaFeAsF, \( d(\text{As-F}) = 2.89 \) Å. For comparison, in LaFeAsO, the separation is \( d(\text{As-O}) = 3.05 \) Å. These differences have the potential to affect \( k_z \) dispersion, that is, the three dimensional character, which is an issue that we study.

B. Calculational Methods

The full-potential linearized augmented plane wave Wien2K package\cite{18} has been used for the electronic structure calculations. We use the Perdew, Burke, and Ernzerhof\cite{19} version of the generalized gradient approximation (GGA) to the exchange-correlation functional within density functional theory. The sphere radii for Ca, H, Fe, As, F, La and Co are taken as 2.50, 1.40, 2.39, 2.12, 2.20, 2.50 and 2.38 Bohr, respectively. The basis set cut-off parameter \( R_{\text{MT}} K_{\text{max}} = 7.0 \) was found to be sufficient. The number of \( \tilde{k} \) points was typically 3000 for the tetragonal unit cell.

The magnetic order commonly observed in these Fe-based two dimensional (2D) materials leads us to study the relative energetics of ordered phases. The primitive tetragonal cell contains two Fe sites, due to the tetrahedral coordination with As ions. The Fe sublattice is square, oriented at 45° with respect to the conven-
tional cell. One possible magnetic order is the Néel type, in which each Fe on its square sublattice is antialigned with its four neighbors; this is \( q = 0 \) AFM order (NAFM). Stripe AFM (SAFM) order corresponds to each Fe atom being antiparallel to its neighbors along the \( \vec{a} \) and \( \vec{b} \) axes, which are second Fe neighbors. SAFM corresponds to \( \vec{Q} = (\pi, \pi, 0) \) order, and typically destroys much of the Fermi surface. The SAFM unit cell size is doubled to \( \sqrt{2} \times \sqrt{2} \times 1 \).

### III. ANALYSIS AND DISCUSSION

#### A. Nonmagnetic phase: a new band

Figure 1 displays the band structure and Fermi surfaces of CaFeAsH in the nonmagnetic phase. There are two cylindrical hole bands around the \( \Gamma=(0,0,0) \) point and two cylindrical electron bands crossing \( E_F \) near the \( M=(\pi,\pi,0) \) point, consistent with those presented by Muraba et al.\(^2\) This band structure is similar to that of the LaFeAsO 1111-type compounds, except for an additional central Fermi surface around the \( Z \) point of the zone. The apparent nesting, which we quantify below, implies that CaFeAsH has a tendency toward the SAFM order phase observed and obtained computationally in other 1111 compounds.

The distinctive feature of these bands is the existence of one band along the \( \Gamma-Z \) direction that shows not only \( k_z \) dispersion but a surprisingly large velocity along \( k_z \). This band gives rise to an additional, ellipsoidal in shape, hole Fermi surface surrounding the \( Z \) point of the zone. Bands near \( E_F \), including the unusual one, have strongly Fe \( d_{xz} \) and \( d_{yz} \) character, with small As 4p character. This new Fermi surface has similarities to one in a magnetically ordered phase of LaFePO that is different from the corresponding Fermi surface of LaFeAsO in the same magnetic structure.\(^2\) The \( k_z \) dispersion is however much less dramatic in LaFePO than in CaFeAsH.

#### B. Electronic structure: H versus F

Muraba et al.\(^3\) have shown that the F 2p and H 1s bands lie well below the Fermi level and are filled, so each is a negatively charged (-1) ion. The densities of states (DOS) in Fig. 2(a) demonstrate that the F states are 2 eV more strongly bound than the H 1s states. Figure 2(a), right panel, also shows the corresponding projected density of states (PDOS) near \( E_F \). These curves likely reflect not the atomic states but the tails of neighboring atomic orbitals that extend into the H and F atomic regions. The point is that these are states near \( E_F \) and their distributions and hybridization are significantly different. It is this and related differences that we pursue in this subsection.

The Fermi surface differences noted in Sec. III.A do not convey the extent of the differences in electronic structure caused by H versus F. Figures 2(b) and 2(c) show the band dispersions, weighted by Fe \( d_{xz} + d_{yz} \) character, in the M-\( \Gamma-X \) and the zone top A-Z-R paths. In CaFeAsF, there are three hole-like bands centered around the \( \Gamma-Z \) line, and the bands display no \( k_z \) dispersion. For CaFeAsH, however, two hole-like bands cross \( E_F \) along the \( M-\Gamma-X \) path while three hole-like bands cross the \( E_F \) along the A-Z-R path. The difference is that in the H compound, a band at \( \Gamma \) 0.2 eV below \( E_F \) in CaFeAsF lies 0.2 eV above \( E_F \) at \( Z \).

Figure 3 shows the bands along the most relevant symmetry lines, weighted this time by the amount of As 4p\(_x\) character. Most prominently, in CaFeAsH, a band is essentially linear from \( \Gamma \) to \( Z \) — very large \( k_z \) dispersion — over an energy span of 1.5 eV. The velocity is \( h v_F = 3.8 \) eV\( \cdot \)\( \AA \), or \( v_F = 5.9 \times 10^6 \) cm/s. The bonding path that enables this large interlayer hopping is Fe-As-H-As-Fe, identified previously.\(^2\) Visible in Fig. 3(b) is the corresponding band in CaFeAsF has a factor of ~7 less dispersion, and in the opposite direction. Also, it lies entirely above \( E_F \) so it does not affect the FS.

A plot through the zone \( Z-\Gamma-Z \) would indicate this dispersive band in CaFeAsH to be Dirac-like at -0.2 eV, except that it (necessarily, by symmetry) becomes...
FIG. 3. (Color online) The contribution of the arsenic atomic 4p_z orbital to (a) CaFeAsH and (b) CaFeAsF bands near E_F. The dashed red lined regions emphasize areas of differences in bands along the Γ-X and Γ-Z directions, see text.

quadratic extremely close to Γ. However, this band disperses downward in both directions in plane, so the point at Γ is a van Hove singularity (vHs) with two normal negative masses m_x = m_y and an extremely small positive mass along the k_z direction, whereupon the band quickly becomes linear (massless). Lying 0.2 eV below E_F this vHs has no influence on low energy properties (including, we expect, superconductivity) but it reflects a highly unusual influence of H in this structure. A very similar linear valence band has been found in the CoAs_3 class of skutterudites where it has been analyzed in detail and found to arise from a cluster orbital of As p states surrounding the open hole in this lattice.

There is an additional H versus F difference evident in Fig. 3(b). Of the pair of hole bands extending from Γ, the Fermi surfaces are circular and nearly identical in CaFeAsH, while the values of k_F in CaFeAsF differ. The cylinders include different numbers of holes, and the larger number of holes in CaFeAsH are balanced by the electrons in the ellipsoid that does not exist in CaFeAsF. This difference also affects the nesting function which can be seen in plots from a careful study that we present below.

C. Fe layer charge: H versus F

Due to the focus of attention in Fe-based pnictides and chalcogenides on Fe orbital occupations and how they relate to properties, we have calculated the difference density ρ(CaFeAsH)−ρ(CaFeAsF), with both calculated at the lattice constant of the former compound. Fig. 4 shows contour plots of this difference density in two planes: a (110) plane containing a FeAs zigzag chain, and a (001) plane through the Fe sites. Positive values designate regions where H attracts charge relative to F.

In the horizontal plane containing Fe sites, Fig. 4(b), H is seen to increase charge in the in-plane orbitals while there is a slight decrease in density in the interstitial region in both planes. In the Fe-As chain direction, Fig. 4(a), the increase in occupation of the Fe in-plane orbitals is again evident. What is more evident is the larger change on As, with charge "transferred" to the p_x, p_y orbitals at the expense of the p_z orbital. In addition, there are changes on the Ca ion of the same size as on Fe. Density is slightly decreased in the interstitial region.

These differences can be quantified by subtracting the orbital occupations, CaFeAsH minus CaFeAsF, within the atomic spheres. In units of 10^{-3}, the differences for Fe are: d_{z^2}, 2.5; d_{xy}, 7.6; d_{x^2-y^2}, 3.5; d_{xz} = d_{yz}, -0.6. The net change for Fe is +12.4×10^{-3} electrons. For As, the analogous changes are: p_x, p_y, 5.9; p_z, -9.9, for a small net change of +1.9×10^{-3} electrons. Thus H induces ~0.015 electrons into the Fe+As spheres compared to F, with the decrease occurring in the interstitial region or in the Ca(H,F) region.

D. Study of H_{1-x}F_x substitution

To determine whether the Fermi surface differences (and band structure differences) are affected by the different lattice parameters we have replaced H by F in CaFeAsH and replaced F by H in CaFeAsF, each with
its own value of the $c$ lattice parameter. The differences in the 3D ellipsoid FS sheet are shown in Fig. 5, replacing H by $H_{1-x}F_x$ in CaFeAsH, for $x=1$, $\frac{1}{2}$, $\frac{1}{3}$, 0, and replacing F by $H_yF_{1-y}$ in CaFeAsF ($y=0$, $\frac{1}{4}$, $\frac{1}{2}$, 1). The effect of $c$ lattice constant is very small; note that in each of the cases $x = 0$, $\frac{1}{2}$, and 1 the surfaces are very similar, with the pinching off of the cylinder occurring around 35-40% H. Evidently the difference is a chemical bonding one that is much more important than intersublayer separations.

**E. Fermi surface nesting**

Fermi surface nesting has played a central role in theories of the Fe-based superconductors, on the one hand accounting for the observed AFM phases in several undoped compounds and on the other as promoting short-range AFM fluctuations that may provide the superconducting pairing mechanism. Both of these aspects focus on the susceptibility peak near $\mathbf{Q} = (\pi, \pi, 0)$, as discussed for Fe pnictides earlier. However, nesting processes, i.e. those which scatter from Fermi surface to Fermi surface, dominate all low energy processes, and it may be important to understand such process throughout the Brillouin zone.

With an eye first toward the magnetic instability in CaFeAsH and CaFeAsF and the wavevector dependence of magnetic fluctuations in the (high temperature) tetragonal structure, we analyze the Fermi surface topology in the nonmagnetic phase. The top view of Fermi surface sheets is shown in Fig. 5(a) and Fig. 5(c), with a copy displaced by $\mathbf{Q}$. The cylindrical, nearly circular, FS sheets at M, and two of the three at $\Gamma$, reflect the 2D character of these Fermi surfaces. As mentioned, a substantial degree of $\mathbf{Q}$ nesting is evident, which we proceed to quantify below. While the circular FSs around $\Gamma$ in CaFeAsH essentially coincide, in the F compound the inner sheet is much smaller, arising from the difference in band structures discussed above.

1. **Formalism**

A quantitative measure of nesting is provided by the Fermi surface nesting function $\xi(\mathbf{q})$, which measures the phase space for scattering through wavevector $\mathbf{q}$ from $\mathbf{k}$ on the FS to $\mathbf{k} + \mathbf{q}$ also on the FS. This function is given by

$$\xi_{\mathbf{q}} = \sum_{\mathbf{k}} \delta(\epsilon_{\mathbf{k}} - \epsilon_F) \delta(\epsilon_{\mathbf{k}+\mathbf{q}} - \epsilon_F)$$

$$= \frac{\Omega_c}{(2\pi)^3} \int_{\mathbf{v}_k \times \mathbf{v}_{k+q}} \frac{dL}_{\mathbf{F}}(\epsilon_F)$$

(1)

where $\Omega_c$ is the cell volume and $\mathbf{v}_k = \nabla_k \epsilon_k$ is the electron velocity ($\hbar=1$). The latter expression, after the volume integral with two $\delta$-function restrictions, gives a geometrical interpretation: it is the integral over the line of intersection $L_{\mathbf{k}}(\epsilon_F)$ of the undisplaced FS and a copy of the FS displaced by $\mathbf{q}$, weighted by the inverse of the cross product of the two velocities. Large contributions arise...
from regions of parallel or antiparallel velocities, \( i.e. \) FS nesting where the velocity cross product becomes small. Contributions are enhanced by small velocities, although that is not an issue with these cylindrical FSs that arise in these compounds, with near constant Fermi velocities around the FSs. \( \xi(\vec{q}) \) has a trivial \( q^{-1} \) divergence at \( q \to 0 \) that is countered by vanishing matrix elements in physical properties involving such long wavelength momentum transfers.

Fermi surface instabilities are more often analyzed in terms of the electronic susceptibility, which is the relevant response function,

\[
\chi_0(q) = \frac{1}{N} \sum_{k,m,n} |M_{km,k+q,m}|^2 \frac{f\left(\varepsilon_{k+q,m}\right) - f\left(\varepsilon_{k,n}\right)}{\varepsilon_{k+q,m} - \varepsilon_{k,n} + i\eta},
\]

(2)

where \( M_{km,k',n} \) is a matrix element of \( \exp(i\vec{q} \cdot \vec{r}) \) between Bloch functions. Though matrix elements often are important, evaluations without matrix elements, corresponding to the “generalized susceptibility,” are much more common. Peaks in the real part of \( \chi_0(q) \) provide the positions of potential wavevector \( \vec{q} \) instabilities. Structure in \( \chi_0(q) \) arises from Fermi surface nesting as well as from other near-\( E_F \) structure in the band structure, viz. the virtual bands that will be discussed below.

Instabilities in the metallic phase, often phrased as instabilities of the Fermi surface, are characterized in terms of peaks in the real part of the electronic susceptibility. The nesting function defined above has its own separate importance. To first (linear) order in small \( \omega \) the connection to the imaginary part \( \chi''_0(q,\omega) \) of the dynamical susceptibility (here matrix elements are set to unity) is

\[
\chi''_0(q,\omega) = \pi\omega\xi(q) + O(\omega^2).
\]

(3)

Thus low energy dissipative processes occur where there is strong weight in \( \xi(q) \), in addition to the Drude interband absorption at small \( |q| \).

2. Discussion of nesting

The behavior of \( \xi(q) \) in both compounds is dominated by the very strong peak precisely at \( \vec{Q} \). In several Fe-based pnictides a maximum at this wave vector also shows up in \( \chi(q,\omega = 0) \) but probably not with nesting as perfect as in this F compound, where Fig. 6(c) illustrated the perfect nesting of the outer electron and hole surfaces in CaFeAsF. The nesting is not as perfect for CaFeAsH, but there are two surfaces of each type of carrier that nest. Note that in the dark regions of these plots there is zero nesting, there are no low energy excitations in these regions, a result of the relatively small Fermi surfaces. Of interest, of course, is that the near perfect nesting at \( \vec{Q} \) indicates a strong tendency for AFM order, and that (electron) doping will degrade the nesting and thereby diminish the tendency for magnetic ordering.

Additional structure can be seen in \( \xi(q) \) that is difficult to see in the susceptibility. In 2D, a circular FS of radius \( k_F \) gives rise to a characteristic divergence \( 1/\sqrt{2k_Fq - q^2} \) as \( q = |\vec{q}| \) approaches \( 2k_F \), then vanishes abruptly for \( q > 2k_F \) where transitions are no longer available. The essentially identical in size circular FSs of CaFeAsF give rise to this characteristic circular peak at radius of \( 2k_F \) around \( |\vec{q}| = 0 \) and almost identically around \( |\vec{q}| = \vec{Q} \), easily visible in the figure. Since phonons acquire renormalization from small (nearly zero) energy processes, \( \xi(q) \) suggests strong phonon renormalization at the nesting wavevector and smaller but definite Kohn anomalies on the \( q \approx 2k_F \) and \( |\vec{q} - \vec{Q}| \approx 2k_F \) circular ridges. Spinwaves should acquire renormalization in the same regions, for the same reasons.

Figure 6(b) provides the nesting function for CaFeAsH in comparison to that of CaFeAsF in Fig. 6(d). In the H compound, the non-circular nature of the ridges of \( \xi(q) \) reflect the not-quite-circular Fermi surfaces. In addition, there is additional structure near the peaks at \( \Gamma \) and \( \vec{Q} \) that arise from the ellipsoidal Fermi surface. Taking the broader picture, however, there is little quantitative difference between the nesting functions of these isoelectronic compounds.

F. Antiferromagnetic phases

Here we discuss magnetic energies and the basic features of the electronic structure of the SAFM phase in undoped CaFeAsH, and draw some parallels with related pnictides. The types of AFM order that are studied are described in Sec. II.B, some of which require a \( \sqrt{2}\times\sqrt{2}\times1 \) supercell (there are two Fe sites within the primitive cell). The energies of four of the simplest magnetic configurations relative to the nonmagnetic phase are presented in Table I, where the Fe atomic sphere moments are also provided. It has been much discussed that Fe moments are overestimated by DFT methods\(^\text{27,34,35}\) so magnetic energies will also be exaggerated, however relative orderings may still be meaningful. These relative energies are in almost the same order as in the related compound CaFeAs\(_2\), but differ somewhat in magnitude\(^\text{27,39}\) reflecting some sensitivity to the Fe-As structure or to the makeup of the blocking layer.

For FM alignment the moment is much smaller and the energy less favorable than for the AFM alignments. The smaller moments account qualitatively for the FM phase being less favorable. The SAFM state is most favorable by a large margin, 120 meV/Fe lower than the BSAFM phase. This lowest energy holds in spite of the SAFM case. The differences in Fe moments in the various phases is evidence that the magnetism has substantial itinerant character. For itinerant magnets the energy gain for magnetism is \( Im^2/4 \), where \( I \) is the Hund’s exchange coupling (also called Stoner \( I \)) and \( m \) is the moment. This behavior is not observed, arising
and M have been mixed and gapped by the SAFM order, as so the cylindrical FSs have vanished, reflecting the massive electronic reconstruction driven by the nesting. The few remaining bands crossing $E_F$ have high velocity.

The resulting electronic structure has noteworthy aspects. The majority states are filled ($d^5$). The minority states seem at first sight to have an expected $e_g-t_{2g}$ splitting, however it is more complex than that. Evidently the SAFM order affects substantially the Fe configuration, based on the loss of tetrahedral symmetry due to magnetic reconstruction of the electronic bonding. There are substantial peaks of minority $d_{x^2-y^2}$ states both at -2 eV and +1.5 eV. The main $d_{xy}$ peak is above $E_F$ but a noticeable fraction lies below $E_F$. $d_{x^2-y^2}$ is mostly above $E_F$; $d_{yz}$ is entirely below $E_F$. The magnetic order thus has a considerable effect on the orbital occupations and thereby the electronic structure.

G. La and Co Doping Phase

Both La and Co provide electronic doping, however Co introduces changes directly into the Fe layer while La simply donates electrons. Here we analyze the differences in electronic structure between indirect (La doping) and direct (Co doping) to identify aspects that may correlate with the question of why $T_c^{max}$ is a factor of two different in the two cases. We present results for the same 25% electron doping of the DOS and different Fermi surfaces of Ca$_{0.75}$La$_{0.25}$FeAsH and Ca$_{0.75}$Co$_{0.25}$AsH in Fig. 8(a). The electron surfaces of course grow in size while hole surfaces decrease, and with substantial changes. We note that the onset $T_c$ in the La-doped case remains around 47 K from $x=0.08$ (metallization) to $x=0.30$, the highest doping level reported. This makes our 25% case representative. The data indicate an unusual insensitivity to doping level, unlike the more common superconducting dome behavior with doping. It may however be possible that there is mesoscopic chemical phase separation in these sample such that the onset $T_c$ represents the value of a single phase. For the homogeneous phases that we are modeling, there are substantial changes with doping.

The projected DOS curves of Fig. 8(b) reflect large differences. For La doping the Fe DOS hardly changes, with $E_F$ increasing rigid-band like. The band filling of course is the same for Co doping. The Co 3$d$ DOS differs from that of Fe, however, with the main peak lying 0.5 eV lower (higher binding energy), with the result that...
FIG. 8. (Color online) (a) The Fermi surfaces and qualitative nesting features of Ca$_{0.75}$La$_{0.25}$FeAsH and CaFe$_{0.75}$Co$_{0.25}$AsH, with electron doping in each case but indirect versus direct respectively. (b) Density of states (DOS) of each atom in Ca$_{0.75}$La$_{0.25}$FeAsH (where only Fe is significant) and CaFe$_{0.75}$Co$_{0.25}$AsH, where Co can be compared with Fe.

The unusual band and extra Fermi surface has been discussed, briefly, in terms of the effect of an incipient band that could help to resolve the differences observed between these two compounds. The differences in orbital occupations in both the Fe and As atomic spheres in the H and F compounds has been quantified. It may be helpful in future to compare these isovalent substitutions with the results on H-doping of the 1111 class, as well. However, with no material-specific theory of $T_c$, the effect on the density and character of states near the Fermi level is substantial, so differences in properties, including the isovalent F compound. Both compounds show very strong Fermi surface nesting at the expected $(\pi, \pi, 0)$ wavevector. The cylindrical Fermi surfaces are different for the two compounds: circular for the F compound but two surfaces with different radii; not so circular for the H compound but nearly identical surfaces.

The 3D ellipsoid pocket remains in each of these cases. Because the velocity along $k_z$ is large, the extent of the ellipsoid along $k_z$ does not increase much, being larger around the waist and larger in the Co doping case. The nesting of electron and hole FSs is severely degraded by 25% doping, a change that is conventionally associated with the destruction of spin density wave order that enables superconductivity to emerge. Although Cooper pairing due to short-range AFM fluctuations that are encouraged by nesting provides the most common explanation for superconductivity in Fe pnictides, there is no longer great interest in details of degraded nesting and we do not provide $\xi(\vec{q})$ for the doped materials. Fig. 8(a) provides some indication of where some nesting remains, by repeating the electron FS cylinders displaced by the most evident displacement. The nesting lies at a larger displacement for Co-doping and displays weaker nesting (the corresponding curvatures of the FSs differ more). Such differences might contribute to the factor of two difference in $T_c$, but the theory is not sufficiently established to pursue this point.

IV. SUMMARY

Density functional based methods have been applied to assess various aspects of the effect of Ca and H, versus La and O, on the underlying electronic structures of these 1111 compounds. Secondarily, the differences in electron doping, by La for Ca or by Co for Fe, have been modeled and analyzed. The isovalent system CaFeAs(H,F) has also been studied. The key feature introduced by H is an unusual band dispersing along $k_z$, as pointed out previously, due to a Fe-As($p_x$)-H-As($p_x$)-Fe hopping path enabled by the hydrogen 1s orbital that is much less strongly bound than the corresponding 2p orbitals in the F compound.

This peculiar band introduces an additional Fermi surface of three dimension character and ellipsoidal shape, something not appearing in other 1111 compounds including the isovalent F compound. Both compounds show very strong Fermi surface nesting at the expected $(\pi, \pi, 0)$ wavevector. The cylindrical Fermi surfaces are different for the two compounds: circular for the F compound but two surfaces with different radii; not so circular for the H compound but nearly identical surfaces. The unusual band and extra Fermi surface has been discussed, briefly, in terms of the effect of an incipient band that could help to resolve the differences observed between these two compounds. The differences in orbital occupations in both the Fe and As atomic spheres in the H and F compounds have been quantified. It may be helpful in future to compare these isovalent substitutions with the results on H-doping of the 1111 class, as in LaFeAsO$_{1-x}$H$_x$ and SmFeAsO$_{1-x}$H$_x$.

The effects of indirect electronic doping (La for Ca) and direct doping (Co for Fe) have been contrasted. The effect on the density and character of states near the Fermi level is substantial, so differences in properties, including $T_c$, are expected. Differences in the nesting of these two electron-doped phases have been described as well. However, with no material-specific theory of $T_c$ in Fe-based superconductors, it would be speculation to try to identify the most important distinctions.

We expect that these several findings will contribute to investigations into the determination of what aspects of these differences are relevant to understand their superconducting distinctions. The observation of a lack of superconducting dome in the phase diagrams of both types of doping, taken at face value, indicates an indifference of superconductivity to details of the electronic structure that is difficult to comprehend, since many studies in Fe-based superconductors point to important effects due to small differences.

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