RHESSI AND SDO/AIA OBSERVATIONS OF THE CHROMOSPHERIC AND CORONAL PLASMA PARAMETERS DURING A SOLAR FLARE

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ABSTRACT

X-ray and extreme ultraviolet (EUV) observations are an important diagnostic of various plasma parameters of the solar atmosphere during solar flares. Soft X-ray and EUV observations often show coronal sources near the top of flaring loops, while hard X-ray emission is mostly observed from chromospheric footpoints. Combining RHESSI with simultaneous Solar Dynamics Observatory/Atmospheric Imaging Assembly (AIA) observations, it is possible for the first time to determine the density, temperature, and emission profile of the solar atmosphere over a wide range of heights during a flare, using two independent methods. Here we analyze a near limb event during the first of three hard X-ray peaks. The emission measure, temperature, and density of the coronal source is found using soft X-ray RHESSI images while the chromospheric density is determined using RHESSI visibility analysis of the hard X-ray footpoints. A regularized inversion technique is applied to AIA images of the flare to find the differential emission measure (DEM). Using DEM maps, we determine the emission and temperature structure of the loop, as well as the density, and compare it with RHESSI results. The soft X-ray and hard X-ray sources are spatially coincident with the top and bottom of the EUV loop, but the bulk of the EUV emission originates from a region without cospatial RHESSI emission. The temperature analysis along the loop indicates that the hottest plasma is found near the coronal loop-top source. The EUV observations suggest that the density in the loop legs increases with increasing height while the temperature remains constant within uncertainties.

Key words: Sun: chromosphere – Sun: flares – Sun: X-rays, gamma rays

Online-only material: color figures

1. INTRODUCTION

Solar flares are one of the most spectacular solar phenomena and are observed over a wide range of the electromagnetic spectrum from radio frequencies to high-energy gamma-ray emission. These solar explosions are associated with the acceleration of large numbers of energetic electrons, plasma heating, and plasma motions. However, the details of plasma heating and particle acceleration and transport in solar flares are still poorly understood. Over the last decade, spatially resolved hard X-ray (HXR) observations with RHESSI (Lin et al. 2002) have notably enriched our understanding of solar flares. Using these X-ray observations the spatial, angular, and energy characteristics of non-thermal electrons in solar flares have been deduced (see Kontar et al. 2011, as a recent review), which resulted in improved understanding of non-thermal electron evolution (see Holman et al. 2011, as a recent review). The detailed analysis of the footpoint emission, the brightest HXR sources in solar flares, revealed the density structure of the chromosphere (Battaglia & Kontar 2011a, 2011b). At the same time, RHESSI provides diagnostics of thermal plasma in the soft-X-ray- (SRX-) emitting loop. This allows for measuring the temperature, emission measure, and density in the corona, namely, the often observed loop-top coronal SXR sources (e.g., Battaglia et al. 2009; Veronig et al. 2005).

However, RHESSI observations are limited to electrons with plasma temperatures higher than around 8 MK. Since the plasma in a flare is unlikely to be isothermal and the temperature is expected to vary over time and for different flare sources one has to extend the observed temperature range to lower values such as observed in EUV. For example, Hinode spectroscopic studies of hot plasmas provided details of plasma motions (e.g., Warren et al. 2008; Liu et al. 2009; Watanabe et al. 2010) and observations of decaying post-flare loops (e.g., Feldman et al. 1995; Varady et al. 2000) provide estimates of the electron density of the bright regions at the loop tops. Line spectroscopy can provide information on the emission measure at different temperatures (e.g., Graham et al. 2011; Raftery et al. 2009), while Krucker et al. (2011) showed how diffraction patterns in TRACE 171Å images can be used to make temperature and density estimates. Since its launch in 2010, the Atmospheric Imaging Assembly (AIA; Lemen et al. 2012) on the Solar Dynamics Observatory (SDO) has been contributing greatly to our knowledge of the solar flare extreme ultraviolet (EUV) emission via full disk imaging, allowing ready comparison with X-ray observations of flaring events. It covers a temperature range from ~5000 K to ~20 MK in nine wavelengths and, even more importantly, provides the spatial resolution allowing to connect the non-thermal and thermal phenomena in solar flares. Combining RHESSI and SDO/AIA observations of flaring loops (complemented by GOES SRX observations) thus allows us to analyze the temperature and density structure from the chromosphere to the corona during a solar flare using two complementary methods.

In this paper we present simultaneous RHESSI, SDO/AIA, and GOES observations taken during the first of three HXR peaks of a limb flare. Fitting RHESSI full Sun spectra with a single temperature model provides the SXR temperature and emission measure while RHESSI images indicate the location of the hot plasma. This is enriched by temperature and emission measure found from GOES, which is most sensitive to temperatures around the RHESSI sensitivity range.
and down to a few MK. HXR footpoint visibility analysis gives information about the chromospheric density structure while SDO/AIA observations are used for differential emission measure (DEM) analysis along the flaring loop. For the DEM analysis a regularized inversion method, developed by Hannah & Kontar (2012), was used. This allows to make full use of the information contained in the six AIA wavelengths sensitive to coronal temperatures as opposed to using filter ratios or single temperature fits.

2. DATA ANALYSIS

We consider a well-observed limb flare that happened on 2011 February 24 with three HXR peaks between 07:29 and 07:33 UT. This GOES M3.5 class flare has one well-defined footpoint during the first HXR peak, the height structure of which was analyzed in some detail by Battaglia & Kontar (2011b). The flare has demonstrated white-light (WL) emission and is also associated with a large filament eruption.3

In the present study we focus on the time of the first HXR peak, around 07:30 UT, because of substantial saturation of AIA images at later times. The spatial structure of the flare as observed by both RHESSI above \( \sim 6 \text{ keV} \) and AIA in the 94 Å, 193 Å, 335 Å, and Helioseismic and Magnetic Imager (HMI; Scherrer et al. 2012) WL 6173 Å continuum channels is shown in Figure 1. The analysis is focused on four physically distinct regions of interest, namely, the coronal source, two parts of the flaring magnetic loop, and the HXR footpoint. Rectangles in Figure 1 show the corresponding regions selected for the analysis.

2.1. Alignment between AIA Images and RHESSI

The co-alignment between AIA images in different wavelength channels is important for the accurate determination of the DEM while correct overlays of AIA images and RHESSI images are necessary to determine cospatial X-ray and EUV sources. The alignment between the different wavelength channels of AIA was determined using the limb-fitting routine aia_coalign_test by Aschwanden et al. (2011). This gives an average co-alignment accuracy of 1.1 ± 0.6 pixels. The co-alignment between RHESSI and AIA images has been checked using the method described in Battaglia & Kontar (2011b). Solar Aspect System and Roll Angle System provide pointing knowledge (Fivian et al. 2002) so that the RHESSI disk center is known to a degree of accuracy better than 0.2′′. The disk center of SDO images can also be accurately verified using AIA full disk images in the chromospheric wavelengths (1600 Å and 1700 Å).

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3 http://www.astro.gla.ac.uk/users/mbattaglia/20110224_online_material/
and WL images from HMI. This cannot be done for the other wavelength channels since the very non-uniform distribution of EUV emission in active regions introduces a strong bias. The SDO position of the solar-disk center is found to deviate by only 0.4/0.6 in the x/y-direction from the RHESSI disk center for the time interval analyzed here. The larger pointing uncertainty is related to the absolute calibration of roll angle of the SDO images (rotation around the disk center). Comparison of similar features and the position of the HXR footpoints relative to the 1700 Å and 1600 Å emission suggests an uncertainty of around 0.2 which corresponds to ∼0.2960°/180° ≈ 3.2′′ uncertainty in the tangential direction near the limb. However, the radial direction uncertainty is less than 0.6 allowing us to have rather precise measurements for near limb events. Since we are investigating regions of interest of several arcseconds in diameter, these uncertainties in relative pointing should not affect the analysis, thus RHESSI images were not shifted nor rotated relative to AIA since any such manipulation would be subjective.4

2.2. AIA Emission Measures and Temperatures Using Regularized Inversion

Full Sun AIA level 1 images were prepped to level 1.5 and normalized by exposure time (see Lemen et al. 2012, and online documentation linked therein). Using only AIA wavelength channels with unsaturated images sensitive to coronal temperatures we find the line-of-sight DEM from AIA data using the method of regularized inversion (Hannah & Kontar 2012). The method allows to reconstruct the DEM, ξ(T), for a selected area or individual pixel,

$$\xi(T) = n^2 \frac{dI}{dT} \text{ (cm}^{-5} \text{ K}^{-1}),$$

where n is the plasma density, T is the temperature, and l is the distance along the line of sight.

Using a set of different wavelength values from a pixel or a region in AIA maps, the method returns a regularized DEM as a function of T. This allows unambiguous determination of the contribution of different temperatures to the total emission measure as opposed to RHESSI, where the spectrum is often consistent with different combinations of EM/T (e.g., Prato et al. 2006). Due to high photon flux during flares, images in some of the wavelengths tend to be saturated during the main phase of the flare so that not all AIA wavelength channels can be used to find the DEM during the peak of the flare. In the present flare the image in the 131 Å wavelength channel was saturated for most of the flaring loop region, while the 193 Å image saturated near the top of the coronal source. This is illustrated in Figure 2 showing the AIA image profiles along the x-direction averaged over 6′′ in the y-direction. This one-dimensional cut across the images (Figure 1) includes the flaring footpoints, the legs of the magnetic loop, and the coronal source.

Therefore, for the subsequent DEM analysis, we use images at 94, 171, 335, 193, and 211 Å for the EUV loop and HXR footpoint regions, omitting 193 Å for the region at the loop top where the image in this wavelength was saturated. We discuss the consequences of this omission in Section 3.1.1. For the inversion we use zeroth-order regularization in a range of temperatures log10 T = 5.7 to log10 T = 7.5 with 16 temperature bins. This corresponds to \(\Delta \log T = 0.11\) in temperature space. The latest version of the AIA response function (including the new CHIANTI fix5) was used.

3. DEM ANALYSIS OF SELECTED SPATIAL REGIONS

We focus the analysis of emission measure, temperature, and density on a few selected, physically meaningful regions.

1. Coronal source, corresponding to the bright SXR source at the top of the loop (region 1 in Figure 1).
2. EUV loop. Two regions are selected at the southern loop leg between the SXR coronal source and the footpoints (regions 2 and 3).
3. Footpoint region. This region is near the origin of the HXR and WL emission but also contains some hot plasma observed in EUV (region 4 in Figure 1).

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4 The SDO images presented in Figure 1 can be rotated to make RHESSI, EUV, and WL images coincide, which is subjective and affects its tangential alignment but not the radial.

5 2012 January, http://sohowww.nascom.nas.nasa.gov/solarsoft/sdo/aia/response/chiantifix_notes.txt
For each of the selected regions we find the total DEM from the region near the first HXR peak of the flare in EUV images taken around 07:30:10 UT. Figure 3 shows the DEM divided by the number of pixels in the area as a function of temperature from each of the four selected regions. The DEM from all of the regions suggests the presence of two temperature components, a low temperature component around 2 MK, and a high temperature component around/above 8 MK.

There are several possible reasons for this bi-model distribution: (a) line-of-sight effect: the low temperature component may not originate from the flaring loop itself, but from higher layers of the atmosphere, (b) true multi-temperature structure of the flaring site, and (c) artifacts related to the SDO/AIA response calibration or the assumed theoretical spectral calculations. To test for (a) we calculated the DEM for the same four regions in images taken around 15 minutes, 10 minutes, 5 minutes, 2 minutes, and 1 minute before the analyzed flare time interval (Figure 4). The time evolution for all four flare regions suggests that the low temperature component is always present, while the high temperature component increases in intensity by more than one order of magnitude during the rise phase of the flare.

The second possibility (b) can be tested by analyzing the pixel-by-pixel DEM parallel as well as perpendicular to the magnetic loop. Such analysis suggests that the intensity of the low temperature component is constant both in parallel and perpendicular direction to the magnetic field, while the intensity of the high temperature component in the parallel direction increases by a factor of three at larger heights consistent with reduced flare emission above the top of the coronal source. This is supported by images in 193 Å that show faint loop structures which are probably not associated with the flare. Thus, it is unlikely that the low temperature component is due to flaring emission but it can be explained as line-of-sight effect. In other words, the high temperature component dominates the actual flare emission and this component is what we focus on.

From the DEM curve in Figure 3 we can then calculate the total emission measure per area. Here we define the total emission measure as the DEM integrated over the main (hot) flare temperature component from \( T_{\text{min}} = 2.5-3 \) MK to \( T_{\text{max}} = 32 \) MK (see Figure 3):

\[
\frac{\text{EM}}{A} = \int_{T_{\text{min}}}^{T_{\text{max}}} \xi(T) dT,
\]

where \( A \) is the area of the analyzed region \( A = (0.6 \times 7.25 \times 10^7)^2 \times N_{\text{pixel}} \) cm\(^2\) with \( N_{\text{pixel}} \) being the pixel number over which the flux was averaged. For comparison with RHESSI and GOES we define a peak, or main, temperature as the first moment of the DEM curve between \( T_{\text{min}} \) and \( T_{\text{max}} \):

\[
\langle T \rangle = \frac{\int_{T_{\text{min}}}^{T_{\text{max}}} \xi(T) \times T dT}{\int_{T_{\text{min}}}^{T_{\text{max}}} \xi(T) dT},
\]

and the second moment gives an uncertainty range of possible temperatures. All temperatures are given in the form \( \langle T \rangle = \)
(\(T\) \(\pm \Delta T\)), where \(T\) is the temperature found from Equation (3) and \(\Delta T\) is the half-width of the second moment. In the next sections we discuss the results for each of the four regions individually. Table 1 summarizes the temperature and emission measure values.

### 3.1. Soft X-Ray Coronal Source

The SXR emission originated from higher altitude than the EUV loop as seen in 94 Å, but coincided with the 193 Å emission. The SXR thermal parameters (temperature and emission measure) of the coronal source were found by fitting the full Sun RHESSI spectrum which is dominated by the coronal source emission at low energies. A thermal component plus a single power-law thick target component with a low energy cutoff were fitted over a time interval of 07:30:00–07:30:40 UT (same time interval as RHESSI imaging). Fits for two different values of the low energy cutoff resulted in comparable goodness of fit, thus we cite the two resulting values of the temperature and emission measure, which also gives a confidence interval for these parameters. The temperature resulted in \(T_{\text{RHESSI}} = 18.5 \pm 1.5\) MK and the emission measure \(EM_{\text{RHESSI}} = (9.4 \pm 3.4) \times 10^{47}\) cm\(^{-3}\), respectively. For comparison with

![Figure 4](image-url)

Figure 4. Pre-flare time evolution of DEMs, using the 94, 171, 211, and 335 Å wavelength channels (top left) and 94, 171, 211, 335, and 193 Å (top right to bottom right; see Section 3.1.1). The black line is the DEM during the main analysis interval, the colored lines give the DEM at different times (from about 15 minutes to 1 minute) before the flare.

(A color version of this figure is available in the online journal.)

| Instrument/Region          | \(T\) (MK) | EM/Area \(\times 10^{47}\) cm\(^{-3}\) | \(n_e\) \(\times 10^{10}\) cm\(^{-3}\) |
|---------------------------|------------|-----------------------------------|----------------------------------|
| RHESSI/coronal source     | 18.5 ± 1.5 | \((1.1 \pm 0.4) \times 10^{47}\)   | \((3.4 \pm 0.6) \times 10^{10}\) |
| AIA/Region 1              | 6.8 ± 1.3  | \((5.8 \pm 1.1) \times 10^{47}\)   | \((8.0 \pm 0.8) \times 10^{10}\) |
| AIA/Region 1 (incl. 193 Å)| 11.6 ± 3.9 | \((2.7 \pm 0.9) \times 10^{47}\)   | \((1.7 \pm 1.0) \times 10^{10}\) |
| AIA/Region 1 (incl. 131 Å)| 12.4 ± 3.0 | \((3.2 \pm 0.8) \times 10^{47}\)   | \((1.9 \pm 0.9) \times 10^{10}\) |
| AIA/Region 1 (incl. 131 and 193 Å) | 13.6 ± 4.0 | \((4.2 \pm 1.1) \times 10^{47}\)   | \((2.2 \pm 1.0) \times 10^{10}\) |
| AIA/Region 2              | 7.5 ± 2.7  | \((1.2 \pm 0.3) \times 10^{47}\)   | \((1.8 \pm 0.2) \times 10^{10}\) |
| AIA/Region 3              | 7.4 ± 2.2  | \((2.1 \pm 0.6) \times 10^{47}\)   | \((1.5 \pm 0.2) \times 10^{10}\) |
| AIA/Region 4              | 8.7 ± 2.0  | \((2.1 \pm 0.8) \times 10^{47}\)   | \((2.4 \pm 0.5) \times 10^{10}\) |
| GOES/coronal source       | 14         | \(3.1 \times 10^{47}\)             |                                  |
AIA emission measures, the emission measure per unit area \( \text{EM}_{\text{RHESSI_4}} \) is used. Here we assume that the RHESSI emission originated from an area the size of the 50% contours in a CLEAN image \( (8.5 \times 10^{17} \text{ cm}^2, \text{CLEAN beam corrected}) \), therefore \( \text{EM}_{\text{RHESSI_4}} = (1.1 \pm 0.4) \times 10^{30} \text{ cm}^{-5} \). Another method is to use filter ratios from the GOES satellites (White et al. 2005), where we assume the same area as for RHESSI. This results in an emission measure per area of \( \text{EM}_{\text{GOES}} \approx 3.1 \times 10^{30} \text{ cm}^{-5} \) at a temperature of 14 MK. The EUV emission measure and temperatures were found as described above and amount to \( 6.8 \pm 1.3 \text{ MK and EM}_A = (5.8 \pm 1.1) \times 10^{28} \text{ cm}^{-5} \).

### 3.1.1. Saturation and High Temperature Wavelength Channels

The 193 Å and 131 Å images were saturated near the SXR coronal source and across the whole loop, respectively, which is why they cannot be used to find the DEM in some regions. However, the images suggest that the 193 Å emission was cospatial with the RHESSI coronal source. Since the 193 Å response has a peak at around 16 MK (Lemen et al. 2012), which is close to the observed RHESSI temperature, it is likely that the two images show the same emission at similar temperature. Thus, by omitting the 193 Å and 131 Å wavelength channels in the DEM calculation, most high temperature emission will be missed. To estimate the effect on the DEM we calculated DEMs by including 193 Å and 131 Å and compared the resulting DEM and total emission measure per area. Figure 5 shows this comparison in the case of exclusion of 193 Å, 131 Å individually and both at a time. When these wavelength channels are included the emphasis is on the high temperature emission, the temperature range extends to the upper limit in the DEM reconstruction (32 MK) and the emission measure per area is found as \( \text{EM}_A = (4.2 \pm 1.1) \times 10^{29} \text{ cm}^{-5} \) which is seven times higher than when these wavelength channels are not used but is still 2.6 times smaller than the RHESSI emission measure (compare Table 1). However, the data values from the 193 Å and 131 Å images are lower limits due to the saturation. This points out the effect of omitting the high temperature response wavelength channels and the obtained coronal source emission measure has to be viewed as a lower limit. Furthermore, the temperature sensitivity range of AIA has to be considered as additional limiting factor. Temperatures around 20 MK are at the upper limit of the AIA sensitivity (the high temperature response of 193 Å peaks at 16 MK). Therefore, if the bulk of the plasma is at such high temperatures it is likely that the DEM analysis will result in a smaller emission measure due to lack of sensitivity, even if all coronal wavelength channels are used.

### 3.2. EUV Loop

The EUV loop was divided into two regions, one just below the SXR coronal source (region 2) and one closer to the HXR footpoint (region 3). The temperature and emission measure in region 2 were \( 7.5 \pm 2.7 \text{ MK and } (1.2 \pm 0.3) \times 10^{29} \text{ cm}^{-5} \), respectively. In region 3, values of \( 7.4 \pm 2.2 \text{ MK and } (2.1 \pm 0.6) \times 10^{29} \text{ cm}^{-5} \) were found. Thus the temperature is constant within the uncertainty along the loop while the emission measure seems to decrease slightly with increasing height. We note that due to line-of-sight effect the emission could partially originate from the loops behind/in front of the loop under study.
3.3. Footpoint

Near the position of the HXR footpoint (region 4) the temperature was 8.7 ± 2.0 MK, while the emission measure per area was found as \((2.1 \pm 0.8) \times 10^{29} \text{ cm}^{-5}\). Note that this region, though partly overlapping with the HXR footpoint emission in images, corresponds to a height of \(\approx 2.2–2.9 \text{ Mm}\). This is above the height of the peak HXR emission (\(\sim 1.6 \text{ Mm}\) at 30–40 keV following Battaglia & Kontar 2011b) and above the height of the WL emission which is cospatial with the HXR emission (Martínez Oliveros et al. 2012). Thus, the emission analyzed here originates from the upper transition region/lower corona rather than the chromosphere (compare Figure 2).

3.4. Comparison of Thermal Parameters and Density

In this section we compare the thermal parameters found from EUV data with the RHESSI observations, as well as with GOES measurements, and we also determine the densities. In a next step, every radial \((x)\)-coordinate is mapped to a height above the photosphere, using the reference coordinate that represents the photospheric height found from visibility analysis of RHESSI footpoints as described in Battaglia & Kontar (2011a, 2011b), where we found a reference position of \(r_0 = 929.4 \pm 0.3\) Mm.

Similar to previous results (e.g., Battaglia et al. 2005; Hannah et al. 2008) the RHESSI temperature is highest, while the GOES emission measure is higher but at a temperature a few MK lower than the RHESSI temperature. The EUV emission measures per area from all four regions of interest are up to two orders of magnitude smaller than the RHESSI and GOES emission measure. This can be explained by a number of reasons. The RHESSI and GOES emission measures are calculated for the full Sun, though in the case of a flare this emission originates predominantly from the flaring site. To compare the emission measure per area with the value from AIA, we assume that the bulk of the RHESSI and GOES emission originates from within an area corresponding to the 50% contours in a RHESSI 6–12 keV image. We can calculate the AIA emission measure of the coronal source for the same area to find \(EM = (1.2 \pm 0.3) \times 10^{29} \text{ cm}^{-5}\). This is a factor of two larger than the emission measure from region 1, but still almost one order of magnitude smaller than the RHESSI emission measure. The main reason, as illustrated in Section 3.1.1, is that emission at highest temperatures is missed by omitting the highest temperature response wavelengths 193 and 131 Å. This is likely worsened by the fact that the temperature at this stage of the flare is higher than the range to which AIA is most sensitive, even if the high temperature response was included, thus the emission measure is smaller.

In a next step we compare the densities found from the different methods. An estimate for the density of the coronal source in SXR can be found by measuring the area \(A_{cs}\) of the coronal source in CLEAN or Pixon images and using \(n_{cs} = \sqrt{EM_{cs}/V_{cs}}\), where the volume is estimated to be \(V_{cs} = A_{cs}^{3/2}\). In this case, using 50% contours in a CLEAN image at 6–12 keV and accounting for the CLEAN beam, we get a volume of \(7.9 \times 10^{26} \text{ cm}^3\) and thus a density of \(n_{cs} = (3.4 \pm 0.6) \times 10^{10} \text{ cm}^{-3}\). In the case of AIA observations the density is given as the square root of the emission measure per area divided by the distance along the line of sight. For the line-of-sight component the observed portion of the loop is assumed to be “cylindrical.” We use the size of the HXR footprint to estimate the loop diameter and approximate the line-of-sight distance to \(5'' = 3.6 \text{ Mm}\) for regions 2–4. The coronal source (region 1) appears extended compared with the loop, therefore a value of 9 Mm is used for the line-of-sight distance. It is often not possible to determine the thermal parameters of the footpoints, even in imaging spectroscopy (Battaglia & Benz 2006; Hudson et al. 1994; McTiernan et al. 1993), either because of the dominance of the coronal source and the dynamic range of RHESSI, or because the plasma is at temperatures to which RHESSI is not sensitive. However, using RHESSI visibility analysis it is possible to determine the chromospheric density near the footpoints by fitting an exponential density model to the energy-dependent positions of the footpoints (Battaglia & Kontar 2011b; Kontar et al. 2010, 2008). Figure 6 gives an overview of the densities as a function of height. The densities in regions 2–4 are all around \(2 \times 10^{10} \text{ cm}^{-3}\). A similar density is found using the SXR RHESSI data. This is consistent with previous densities found from RHESSI thermal analysis and indicates a constant density along the loop. For these calculations a filling factor of one was assumed. It cannot be excluded that the filling factor is smaller. However, recent analysis of several flare loops suggests a value that is generally close to one (Guo et al. 2012a, 2012b). Even a filling factor as low as 0.1 would lead to a density only three times higher which is still feasible for a flaring loop.

Since the density found from HXR observations is the neutral plasma density, and AIA observes hot, ionized plasma, combining these observations gives a complete picture of the density in the solar atmosphere during a flare. The combined observations represent a steeply decreasing density from the photosphere to a height of about 4″–5″ with a flatter decrease or near constant value at larger heights, similar to models such as from Vernazza et al. (1981) and Avrett & Loeser (2008). The Avrett & Loeser (2008) quiet-Sun model is shown in Figure 6 along with a double-exponential density model with two scale heights (220 km at small heights, 17,000 km at larger heights) that follows the observed pattern. The AIA-determined density of region 1 is lower by about a factor of four compared with RHESSI because of the small emission measure.

4. DISCUSSION AND CONCLUSION

In this work we present a comprehensive study of the thermal parameters (emission measure, temperature) and the density in four different regions of a flaring loop, using both X-ray as well as EUV data. Detailed temperature, emission measure, and density analysis along the loop near the flare HXR peak indicate that the hottest plasma is found near the coronal loop-top source, in line with previous observations. The EUV observations suggest that the density in the loop legs increases with decreasing height, but the temperature is constant within uncertainties. The footpoint structure is such that the plasma is neutral or partly ionized up to \(\sim 1–2 \text{ Mm}\) with a steep, unresolved rise to \(\sim 5 \text{ Mm}\) 1–2 Mm above the height of the HXR and WL emission.

The DEM in EUV was found using the method of regularized inversion developed by Hannah & Kontar (2012). We show that the method is capable of recovering two temperature components, one at around 2 MK and one at or around 8 MK. The low temperature component around 2 MK is observed in all four regions and is only weakly changing with time, while the intensity of the high temperature component increases by more than one order of magnitude at the time of the flare. Therefore, the low temperature component can be attributed to line-of-sight effect. The existence of this component has to be kept in mind when DEMs are determined by fitting single
Figure 6. Top: emission measure per area as a function of height found for the different regions and using different methods. Red: region 4; green and blue: regions 3 and 2; turquoise: coronal source from AIA (region 1); purple: coronal source from RHESSI. Middle: temperature. Bottom: densities. The black lines give the footpoint density model and its uncertainties as found via HXR visibility analysis (taken from Battaglia & Kontar 2011b). The dashed line gives the sum of two exponential density functions with scale heights of 220 km and 17000 km, respectively. The dotted line indicates the quiet-Sun model by Avrett & Loeser (2008). (A color version of this figure is available in the online journal.)

The loops visible in 94 Å are around 5" lower than the 193 Å loop structure. While the brightest SXR feature seen by RHESSI in 10–15 keV is undoubtedly the coronal source, the brightest 193 Å source is in fact the region to the north of the flare located near \( x = -925", y = 285" \). This is intriguing since the emission seems to be part of the active region and as such contains a significant fraction of the thermal energy. However, this region is not associated with either RHESSI source at any stage of the flare.

The saturation in the “hottest” wavelength channels 193 Å and 131 Å disallows the study of the hottest plasma and limits the temperature range of which plasma can be analyzed.
Comparison of the total AIA emission measure with RHESSI and GOES emission measures suggests that the EUV emission is up to more than one order of magnitude smaller than the X-ray emission measure. This can be partly attributed to omitting the highest temperature wavelength channels in the analysis due to saturation, and it emphasizes the importance of using all six “coronal” wavelength channels for the DEM reconstruction. An additional likely cause is that the bulk of the plasma is at a temperature higher than the main sensitivity range of AIA. Cooler events (i.e., with a RHESSI temperature closer to 10 MK) are probably better suited for more detailed comparative studies. In addition, in weaker events the likelihood for AIA saturation is reduced.

Combined RHESSI/SDO-AIA studies allow to investigate in detail a broad temperature range from 1–2 MK up to ∼20 MK. Moreover, the use of DEM analysis helps to quantify the temperature, emission, and density structure in the flaring loop and fills the gap between the coronal source and footpoints that are observed in SXR and HXR.

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REFERENCES

Aschwanden, M. J., Boerner, P., Schrijver, C. J., & Malanushenko, A. 2011, Sol. Phys., 364
Avrett, E. H., & Loeser, R. 2008, ApJS, 175, 229

Battaglia, M., & Benz, A. O. 2006, A&A, 456, 751
Battaglia, M., Fletcher, L., & Benz, A. O. 2009, A&A, 498, 891
Battaglia, M., Grigis, P. C., & Benz, A. O. 2005, A&A, 439, 737
Battaglia, M., & Kontar, E. P. 2011a, ApJ, 735, 42
Battaglia, M., & Kontar, E. P. 2011b, A&A, 533, L2
Feldman, U., Seely, J. F., Doschek, G. A., et al. 1995, ApJ, 446, 860
Fivian, M., Hemmeck, R., McHedlishvili, A., & Zehnder, A. 2002, Sol. Phys., 210, 87
Graham, D. R., Fletcher, L., & Hannah, I. G. 2011, A&A, 532, A27
Guo, J., Enslie, A. G., Kontar, E. P., et al. 2012a, A&A, 543, A53
Guo, J., Enslie, A. G., Massone, A. M., & Piana, M. 2012b, ApJ, 755, 32
Hannah, I. G., Christe, S., Krucker, S., et al. 2008, ApJ, 677, 704
Hannah, I. G., & Kontar, E. P. 2012, A&A, 539, A146
Holman, G. D., Aschwanden, M. J., Aurass, H., et al. 2011, Space Sci. Rev., 159, 107
Hudson, H. S., Strong, K. T., Dennis, B. R., et al. 1994, ApJ, 442, L25
Kontar, E. P., Brown, J. C., Enslie, A. G., et al. 2011, Space Sci. Rev., 159, 301
Kontar, E. P., Hannah, I. G., Jeffrey, N. L. S., & Battaglia, M. 2010, ApJ, 717, 250
Kontar, E. P., Hannah, I. G., & MacKinnon, A. L. 2008, A&A, 489, L57
Krucker, S., Raftery, C. L., & Hudson, H. S. 2011, ApJ, 734, 34
Lemen, J. R., Title, A. M., Akin, D. J., et al. 2012, Sol. Phys., 275, 17
Lin, R. P., Dennis, B. R., Hurford, G. J., et al. 2002, Sol. Phys., 210, 3
Liu, R., Wang, H., & Alexander, D. 2009, ApJ, 696, 121
Martinez Oliveros, J.-C., Hudson, H. S., Hurford, G. J., et al. 2012, ApJ, 753, L26
McTiernan, J. M., Kane, S. R., Loran, J. M., et al. 1993, ApJ, 416, L91
Prato, M., Piana, M., Brown, J. C., et al. 2006, Sol. Phys., 237, 61
Raftery, C. L., Gallagher, P. T., Milligan, R. O., & Klimchuk, J. A. 2009, A&A, 494, 1127
Scherrer, P. H., Schou, J., Bush, R. I., et al. 2012, Sol. Phys., 275, 207
Varady, M., Fludra, A., & Heinzel, P. 2000, A&A, 355, 769
Vernazza, J. E., Avrett, E. H., & Loeser, R. 1981, ApJS, 45, 635
Veronig, A. M., Brown, J. C., Dennis, B. R., et al. 2005, ApJ, 621, 482
Warren, H. P., Feldman, U., & Brown, C. M. 2008, ApJ, 685, 1277
Watanabe, T., Hara, H., Sterling, A. C., & Harra, L. K. 2010, ApJ, 719, 213
White, S. M., Thomas, R. J., & Schwartz, R. A. 2005, Sol. Phys., 227, 231