Quark transverse spin-orbit correlations

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We extend the study of quark spin-orbit correlations in the nucleon to the case of transverse polarization. At the leading-twist level, this completes the spin structure of the quark kinetic energy-momentum tensor. In particular, we revisit the transversity decomposition of angular momentum proposed a decade ago by Burkardt and introduce a new transverse correlation, namely between quark transversity and orbital angular momentum. We also provide for the first time the Wandzura-Wilczek expression for the second Mellin moment of twist-3 transversity generalized parton distributions, along with a new sum rule. Based on lattice calculation results, we conclude that the quark transverse spin-orbit correlation is negative for both up and down flavors, just like in the longitudinal case.

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I. INTRODUCTION

Understanding the nucleon spin structure is one of the key questions in hadronic physics. It opens a window on a wide range of non-perturbative effects in quantum chromodynamics (QCD) currently studied at many facilities such as Jefferson Lab, RHIC and COMPASS [1], and is a major pillar of the physics case of the future electron-ion collider (EIC) [2]. Although the proper decomposition of the nucleon spin into quark and gluon contributions constitutes one of the fundamental motivations in this field, see e.g. [3, 4], the spin structure turns out to be much richer owing to spin-orbit correlations [6–8].

In a former paper [7], the quark longitudinal spin-orbit correlation was studied in detail by performing a (chiral-even) helicity decomposition of the quark energy-momentum tensor. It has, in particular, been shown that the quark longitudinal spin-orbit correlation can quantitatively be expressed in terms of parton distributions. Both current phenomenological extractions based on experimental data and lattice calculations indicate that the quark spin is, in average, opposite to the quark kinetic orbital angular momentum (OAM).

In this Letter, we discuss the quark transverse spin-orbit correlation by revisiting the (chiral-odd) transversity decomposition of the quark energy-momentum tensor considered a decade ago by Burkardt [9, 10]. Mimicking the approach used by Ji to relate angular momentum contributions to generalized parton distributions (GPDs) [11], Burkardt decomposed the symmetric energy-momentum tensor and introduced accordingly the correlation between quark transversity and total angular momentum. Here we consider the more general asymmetric energy-momentum tensor leading to another transverse correlation, now between quark transversity and OAM.

The Letter is organized as follows: In section II we define the quark transverse spin-orbit correlation operator and express the corresponding expectation value in terms of tensor generalized form factors. In section III we relate these generalized form factors to moments of measurable parton distributions and derive for the first time the Wandzura-Wilczek expression for the second Mellin moment of twist-3 transversity generalized parton distributions, along with a new sum rule. In section IV we compare the various contributions obtained on the lattice with relativistic quark model predictions, provide an estimate of the quark transverse spin-orbit correlation, and we conclude the paper with section V.

II. QUARK SPIN-ORBIT CORRELATIONS

A. Decomposition based on polarization

It is well known that the quark field operator can be decomposed into right- and left-handed contributions

\[ \psi = \psi_R + \psi_L, \quad \psi_{R,L} = \frac{1}{2}(1 \pm \gamma_5)\psi. \]

The quark number and helicity light-front operators can then respectively be seen as the sum and difference

\[ \int d^3x \bar{\psi}_R \gamma^+ \psi = \hat{N}_R^q + \hat{N}_L^q, \]
\[ \int d^3x \bar{\psi}_L \gamma^+ \gamma_5 \psi = \hat{N}_R^q - \hat{N}_L^q \]

of the right and left-handed densities

\[ \hat{N}_{R,L}^q = \int d^3x \bar{\psi}_{R,L} \gamma^+ \psi_{R,L}. \]
where $a^\pm = \frac{1}{2}(a^0 \pm a^3)$ for a generic four-vector $a$, and $d^3x = dx^0 dx^1 dx^2 dx^3$.

Alternatively, the quark field operator can be decomposed into up and down transverse polarizations[12]

$$\psi = \psi_\uparrow + \psi_\downarrow, \quad \psi_\uparrow = \frac{1}{2}(1 \pm \gamma^5)\psi$$

with $j = 1$ or 2. While the sum of up and down densities naturally gives the quark number operator, their difference defines the so-called quark transversity

$$\int d^3x \bar{\psi}(x)\gamma^\perp(x)\psi(x) = \hat{N}_\uparrow - \hat{N}_\downarrow, \quad \hat{N}_\uparrow, \hat{N}_\downarrow$$

where

$$\hat{N}_\uparrow, \downarrow = \int d^3x \bar{\psi}(x)^+ \gamma^\perp(x)\psi(x)$$

The same decompositions can be performed with the quark light-front OAM operator

$$\sqrt{2} \epsilon_{\alpha\beta} \int d^3x \bar{\psi}(x)^+ \gamma^\perp(x)\gamma^\perp(x)\frac{\partial\hat{a}}{\partial x^\alpha} \psi(x) = \hat{L}_R^{\alpha\beta} + \hat{L}_L^{\alpha\beta}$$

where

$$\hat{L}_a^{\alpha\beta} = \sqrt{2} \epsilon_{\alpha\beta} \int d^3x \bar{\psi}(x)^+ \gamma^\perp(x)\gamma^\perp(x)\frac{\partial\hat{a}}{\partial x^\alpha} \psi(x)$$

with the convention $\epsilon_{0123} = +1, a = R, L, \uparrow, \downarrow$.

B. Parametrization

We find that the non-forward matrix elements of $\hat{T}^{\lambda\mu\nu}_{q5}$ can be parametrized in terms of seven generalized form factors (GFFs)

$$\langle p', s'|\hat{T}^{\lambda\mu\nu}_{q5}(0)|p, s \rangle = \mathcal{M}(p', s')\Gamma^{\lambda\mu\nu}_{q5}(p, s)$$

with

$$\Gamma^{\lambda\mu\nu}_{q5}(p, s) = \frac{\rho^{\lambda\mu\nu}_{q5}}{2Mx_5} A_T^{\lambda\mu\nu}(p, t) + \frac{\rho^{\lambda\mu\nu}_{q5}}{2Mx_5} A_T^{\lambda\mu\nu}(p, t)$$

where $s$ and $s'$ are the initial and final rest-frame polarization unit vectors, $M$ is the nucleon mass, $p = \frac{p_T}{m}$ is the average four-momentum, and $t = \Delta^2$ is the square of the four-momentum transfer $\Delta = p' - p$. Note that the last term is totally antisymmetric over all three Lorentz indices, so that $\rho^{\lambda\mu\nu}_{q5} = \rho^{\lambda\mu\nu}_{q5}$.

To recover the twist-2 parametrization of Hager and Diehl[13, 14], one has to symmetrize over the pair of indices $(\lambda\mu)$, and remove all the traces[13]. As a result, the tilde GFFs become redundant

$$3A_T^{\lambda\mu\nu}(t) \rightarrow -D_T^{\lambda\mu\nu}(t),$$

without summation over $j$ in[12]. These are the diagonal components of a $3 \times 3$ matrix whose entries are the directions of quark polarization and OAM.

The longitudinal spin-orbit correlation[11] has been studied in[12]. In this Letter, we focus on the transverse spin-orbit correlation[12] which can conveniently be rewritten as (once again without summation over $j$)

$$\hat{C}_j^q = \sqrt{2} \epsilon_{ji} \int d^3x \bar{\psi}(x)^+ \gamma_5 x^i \frac{\partial\hat{T}^{\lambda\mu\nu}_{q5}}{\partial x^j} \psi(x)$$

with $\hat{T}^{\lambda\mu\nu}_{q5}$ the quark energy-momentum tensor where $\gamma^\mu$ has been replaced by $i\sigma^{\lambda\mu\nu}\gamma_5$

$$\hat{T}^{\lambda\mu\nu}_{q5}(x) = \bar{\psi}(x)i\sigma^{\lambda\mu\nu}\gamma_5 \frac{\partial\hat{D}^\nu}{\partial x^j} \psi(x).$$
\[ s' = s = (s_1, s_2) \), we obtain
\[
\langle p', s | \tilde{T}^{\lambda \nu}_{q_{s}} | p, s \rangle = \left[ \frac{2 P^+}{M} \bar{P}_{\lambda} + \frac{M P^+}{M} \right] (B^T t - D_2 t) \\
+ \frac{1}{2} \int \frac{dz}{2\pi} e^{i x t + z} \langle p', s | \bar{\psi}(-\frac{e}{2}) i \sigma^a \gamma_{\lambda} \psi(\frac{e}{2}) | p, s \rangle = \frac{i z t}{2\pi} \tau(p', s) \Gamma_{q_{\lambda}}^{a} u(p, s)
\]
with \( W = \mathcal{P} \exp[ig \bar{D}_{\lambda} T(x, \xi, t) + \frac{\gamma^\lambda - \Delta^\lambda}{2M}] E_{T}^q(x, \xi, t) \]
written in the symmetric frame \( P_\perp = 0 \).

### A. Equations of motion

The relations for the tilde GFFs can be obtained using the following QCD identities

\[
\bar{\psi} i \gamma^\lambda \gamma_5 \bar{D}_\mu \psi = 2m \bar{\psi} \gamma^\lambda \gamma_5 \psi + i \partial^\lambda (\bar{\psi} \gamma_5 \psi),
\]
\[
\bar{\psi} i \gamma^\lambda \gamma_5 \bar{D}_\mu \psi = -2\imath \epsilon^{\mu \nu \alpha \beta} \partial_\alpha (\bar{\psi} \gamma_5 \psi),
\]
where \( m \) is the quark mass. Taking the corresponding matrix elements and using some Gordon and \( e \)-identities, we find

\[
\langle p', \bar{s} | \bar{\psi} \gamma^\lambda \gamma_5 \psi | p, s \rangle = \pi(p', s') \Gamma_{q_{\lambda}} \psi u(p, s),
\]
\[
\langle p', \bar{s} | \bar{\psi} \gamma^\mu \gamma_5 \psi | p, s \rangle = \bar{\pi}(p', s') \Gamma_{q_{\mu}} \psi u(p, s)
\]
with

\[
\Gamma_{q_S} = \Sigma_{\tau}(t),
\]
\[
\Gamma_{q_P} = \gamma_5 \Pi_{\tau}(t),
\]
\[
\Gamma_{q_{\lambda}}^{a} = \gamma^\mu \gamma_5 G_A^a(t) + \frac{\Delta^\mu}{2M} G_5^a(t).
\]

The quark transverse spin-orbit correlation is therefore given by the expression

\[
\tilde{E}_T^q(x, \xi, t) = \frac{\epsilon^\lambda}{\epsilon^\lambda} \left( \frac{1}{2} \int \frac{dx}{2\pi} \bar{T}_{q_{s}}^{\lambda \nu}(x, 0, 0) + \frac{1}{2} \tilde{E}_T^q(x, 0, 0) \right)
\]

\[
- \frac{2}{3} \Sigma_{q_{\tau}}(0) - \frac{\epsilon^\lambda}{\epsilon^\lambda} C_A^q(0),
\]

### III. LINK WITH PARTON DISTRIBUTIONS

No fundamental probe coupling to \( T_{q_{s}}^{\lambda \nu} \) is known in particle physics. It is however possible to relate the corresponding GFFs to specific moments of measurable parton distributions. From the leading-twist component \( T_{q_{s}}^{\lambda \nu} \), we find in agreement with [13, 14]

\[
\int dx xH_{T}^q(x, \xi, t) = -\frac{1}{2} A_{T}^q(t) - B_{T}^q(t) + D_{2}^q(t) - A_{T}(t),
\]
\[
\int dx xE_{T}^q(x, \xi, t) = A_{T}^q(t) + B_{T}^q(t) = B_{T}(t),
\]
\[
\int dx x\tilde{H}_{T}^q(x, \xi, t) = -\frac{1}{2} A_{T}(t) = \tilde{A}_{T}(t),
\]
\[
\int dx x\tilde{E}_{T}^q(x, \xi, t) = -\xi C_{T}^q(t) = -2\xi \tilde{B}_{T}(t),
\]
where the skewness variable is given by \( \xi = -\Delta^+/2P^+ \) and the functions \( H_{T}^q(x, \xi, t) \), \( E_{T}^q(x, \xi, t) \), \( \tilde{H}_{T}^q(x, \xi, t) \), and \( \tilde{E}_{T}^q(x, \xi, t) \) are the GPDs parametrizing the non-local twist-2 tensor light-front quark correlator [18, 20]
where $\tilde{E}_T^q(x, \xi, t) \equiv 2\tilde{H}_T^q(x, \xi, t) + E_T^q(x, \xi, t)$. Interestingly, it is very similar to the corresponding expression for the longitudinal spin-orbit correlation [7]

$$C_T^q = \frac{1}{2} \int dx x \tilde{H}_q(x, 0, 0) - \frac{1}{2} [F_1^q(0) - \frac{m}{2\pi} H_1^q(0)]$$

and Ji’s relation [11] for the quark OAM

$$L_T^q = \frac{1}{2} \int dx x [H_q(x, 0, 0) + E_q(x, 0, 0)] - \frac{1}{2} G_A^q(0).$$

One might be surprised that Eq. (43) involves thirds instead of halves. They appear because of the factors 3 in Eqs. [33]-[36], which trace back to the fact that $C_T^q$ is defined from a rank-3 tensor, while $L_T^q$ and $C_T^q$ are defined from rank-2 tensors.

Let us stress that the quark transverse spin-orbit correlation introduced in this Letter corresponds actually to the correlation between quark transversity and OAM, $(L_T^q T_2)$. The similarity of our result [33] with Eq. [13], which can be understood as the difference between total angular momentum and spin $(L_T^q S_T^N) = \langle J_T^q S_T^N \rangle - \langle S_T^N S_T^N \rangle$ according to [8] [11], hints to the identifications $(J_T T_2) \propto \int dx x [H_2^T(0, 0) + \frac{1}{2} E_T^q(x, 0, 0)]$ and $(S_T T_2) \propto \Sigma_q(0)$ in the chiral limit $m = 0$. In particular, it suggests that the scalar charge can be interpreted as a measure of the correlation between quark transversity and spin. This interpretation can be further supported by the following simple reasoning in instant form. The difference between spin $\psi^j \sigma^{ij} \psi$ and transversity $\psi^j \langle x^i \rangle \sigma^{ij} \psi$ is a factor $\gamma^0$, and hence of relativistic nature [21]. The correlation between spin and transversity then reads $\psi^j \langle x^i \rangle \sigma^{ij} \psi$ (without summation over $i, j$), which simplifies to the scalar bilinear $\psi^j \gamma^0 \psi = \overline{\psi} \gamma^0 \psi$.

### B. Twist-3 tensor GPDs

The tilde GFFs can alternatively be expressed in terms of twist-3 tensor GPDs. From the twist-3 components $\tilde{T}^{j+}_{q^0}$ and $\tilde{T}^{j-}_{q^0}$, we obtain the following relations

$$\int dx x H_T^q(x, 0, 0) = -\xi \left[ \tau C_T^q(t) + D_T^q(t) + \tilde{D}_T^q(t) \right],$$

$$\int dx x E_T^q(x, 0, 0) = \xi \left[ \tau C_T^q(t) + D_T^q(t) + \tilde{D}_T^q(t) \right],$$

$$\int dx x \tilde{H}_T^q(x, 0, 0) = (1 - \tau) B_T^q(t) + \tilde{B}_T^q(t) + D_T^q(t),$$

$$\int dx x \tilde{E}_T^q(x, 0, 0) = \xi \left[ (1 - \tau) A_T^q(t) + \tilde{A}_T^q(t) + D_T^q(t) \right],$$

where the functions $H_T^q(x, \xi, t)$, $E_T^q(x, \xi, t)$, $\tilde{H}_T^q(x, \xi, t)$, and $\tilde{E}_T^q(x, \xi, t)$ are the GPDs parametrizing the non-local twist-3 tensor light-front quark correlators [20]

$$\frac{1}{2} \int \frac{dz}{2\pi} e^{ixe^{z+}z} \langle p', s' | \bar{\psi}(z^+ \tau) \sigma^{ij} \gamma_5 W \psi(z^+ \tau) | p, s \rangle = \frac{M}{2\sqrt{p'^+ + M}} \langle p', s' \rangle T^{j+}_{q^{-}} u(p, s),$$

$$\frac{1}{2} \int \frac{dz}{2\pi} e^{ixe^{z+}z} \langle p', s' | \bar{\psi}(z^+ \tau) \sigma^{ij} + \gamma_5 W \psi(z^+ \tau) | p, s \rangle = \frac{M}{2\sqrt{p'^+ + M}} \langle p', s' \rangle T^{j+}_{q^{+}} u(p, s).$$

with

$$\Gamma^{j+}_{q^{+}} = -ie_T^{j+} \left[ \gamma^0 H_2^q(x, \xi, t) + \frac{ie_T^{j+}}{2\pi M} E_2^q(x, \xi, t) \right],$$

$$\Gamma^{j+}_{q^{-}} = \gamma^0 \gamma_5 H_2^q(x, \xi, t) + \frac{p^+ + m}{M} E_2^q(x, \xi, t).$$

Since the eight (twist-2 and twist-3) tensor GPD moments are expressed in terms of seven GFFs, there exists a sum rule among them. Adding Eqs. (41) and (42) and using Eq. (29), we find

$$\int dx x \left[ (1 - \tau) \tilde{E}_T^q + H_2^q + E_2^q \right] = 0.$$

Moreover, using the relations [33]-[36], we obtain

$$\frac{1}{\tau - 1} \frac{1}{\xi} \int dx x \left[ H_2^q + \tau E_2^q \right] = \frac{1}{\xi} \Sigma_q$$

$$\int dx x \tilde{H}_2^q = \frac{m}{3M} G_A^q$$

$$\frac{1}{\xi} \int dx x \tilde{E}_2^q = -\frac{m}{3M} G_T^q + \frac{1}{\xi} \Pi_T$$

In the massless quark limit, these expressions provide the Wandzura-Wilczek approximation to the second Mellin moment of twist-3 tensor GPDs. Note that they are exact since no three-parton functions were involved in the derivation, similar to what was observed in the chiral-even sector [22, 24]. Thanks to these results, the quark transverse spin-orbit correlation can now be written in terms of tensor GPDs only

$$\frac{\gamma^0}{\xi} \frac{p^+}{M} C_T^q = \int dx x H_2^q(x, 0, 0) + \frac{1}{\xi} E_2^q(x, 0, 0)$$

$$+ \int dx x \tilde{H}_2^q(x, 0, 0) + 2H_2^q(x, 0, 0),$$

where $H_2^q(x, 0, 0) = \lim_{\xi \rightarrow 0} \frac{1}{\xi} H_2^q(x, \xi, t)$. It is the chiral-odd analogue of the Penttinen-Polyakov-Shuvaev-Strikan relation for the Ji or kinetic OAM [22, 24]}

$$L_T^q = -\int dx x G_2^q(x, 0, 0)$$
and of the relation we found in [3] for the quark longitudinal spin-orbit correlation
\[
C_s^q = -\int dx [\tilde{G}_s^q(x, 0, 0) + 2\tilde{G}_s^{1(q)}(x, 0, 0)].
\] (60)

Note that this time both twist-2 and twist-3 GPDs are necessary to express the quark transverse spin-orbit correlation. This may be due to the fact that transversity does not coincide with transverse spin \[12\].

IV. DISCUSSION

A. Burkardt’s correlation

In the former sections, we worked with the asymmetric quark kinetic energy-momentum tensor and performed a decomposition in terms of quark transversity states. The quark transverse spin-orbit correlation \[C_s^q = \langle L_s^2 T_s^2 \rangle\] can therefore alternatively be seen as the transversity asymmetry of the quark OAM, i.e., \[\langle \delta^s L_s^q \rangle\] following Burkardt’s notation.

This has to be contrasted with the work of Burkardt in [8, 10] which is based on the Belinfante or symmetric quark kinetic energy-momentum tensor [25, 27]. Since in this case the total angular momentum assumes a purely orbital form, Burkardt interpreted his correlation as the transversity asymmetry of the quark total angular momentum \[\langle \delta^s J_s^q \rangle\]. It may be tempting to identify it with the transversity correlation between quark transversity and total angular momentum \[\langle T_s^q \rangle\], just like we identified the quark transverse spin-orbit correlation \[C_s^q = \langle L_s^2 T_s^2 \rangle\] with the transversity asymmetry of the quark OAM \[\langle \delta^s L_s^q \rangle\]. This is, however, not consistent since \[\langle T_s^q \rangle \neq \langle \frac{1}{2} \hat{T}_s^{(-1)} \rangle\] as one can see from the QCD identity \[8, 33\].

In the light-front formalism, Burkardt’s quark operator is given by
\[
\tilde{C}_s^q(x) = \sqrt{2} \int d^3x \left[ x^\mu \hat{T}_q^{1(\nu^2)} - x^\nu \hat{T}_q^{1(\nu^2)} \right]
\] (61)
which is like our operator \[\tilde{C}_s^q\] time-dependent \[38, 39\]. Symmetrizing the expansion \[39\] over the pair of indices \(\mu\nu\), we find
\[
\langle \delta^s J_s^q \rangle \equiv \frac{\langle P_s \hat{C}_s^q(-P_s) \rangle}{\langle P_s \hat{P}_s \rangle} = -\frac{M}{2\sqrt{2}F_p} \left( B_s^q + 2\tilde{B}_s^q - 2D_s^q \right).
\] (62)

Note that the GFF \[\tilde{D}_s^q\] naturally drops out of the final result since it is associated with a Lorentz structure antisymmetric in the pair of indices \(\mu\nu\).

Actually, Burkardt used the instant-form (IF) formalism, where the quark operator is defined as
\[
\tilde{C}_s^q,_{IF} = \int d^3x \left( x^2 \hat{T}_q^{1(03)} - x^3 \hat{T}_q^{1(02)} \right).
\] (63)

We obtain in this case (once again with \[P_s = 0\])
\[
\langle \delta^s J_s^q \rangle_{IF} = \frac{1}{2} \left( \frac{E - M}{M} B_s^q + D_s^q \right)
\] (64)
which is reminiscent of Leader’s result for the transverse Belinfante angular momentum \[3, 28\]
\[
\langle J_s^q \rangle_{IF} = \frac{1}{2} \left( \frac{E - M}{M} B_s^q + (A^q + B^q) \right).
\] (65)

In the rest frame, we recover Burkardt’s result
\[
\langle \delta^s J_s^q \rangle_{IF, rest} = \frac{1}{2} D_s^q = \frac{1}{2} (A_{T20} + 2\tilde{A}_{T20} + B_{T20}),
\] (66)
where we have used Eqs. \[21\] and \[23\].

At first sight, it may seem odd that the light-front and instant-form results \[22\] and \[24\] have different high-energy limits. This is because the transverse OAM light-front operator involves the \(a^-\) component, whereas the instant-form operator involves \(a^+ = \frac{1}{2}(a^+ - a^\perp)\). Therefore, in the high-energy limit \(E \gg M\), the light-front result behaves as \(O(E\^{-1})\) whereas the instant-form result behaves as \(O(E)\). In other words, the instant-form operator contains contributions which are of higher-twist compared to the corresponding light-front operator.

Beside the instant-form approach, Burkardt proposed in [8] a heuristic derivation of Eq. \[66\] based on the light-front operator \[\hat{T}_q^{j(\pm)}\], allowing for an intuitive partonic interpretation in impact-parameter space. Considering the matrix element of the operator \(\sqrt{2} \int d^3x x^3 \hat{T}_q^{j(\pm)}\), one obtains \(\frac{\mu^+ M}{\sqrt{2}M} B_s^q\) which coincides, as expected, with the instant form result \[66\] in the infinite-momentum frame.

Working in the rest frame to invoke rotational symmetry, Burkardt added an extra term \(\frac{1}{2} \int dx x H_t^q(x, 0, 0) = \frac{1}{2}(D_s^q - B_s^q)\) to account for an overall transverse displacement of the center of light-front momentum with respect to the origin, a relativistic effect associated with rotating bodies. Although quite appealing, this interpretation is however not satisfactory since, as stressed in \[8, 22, 31\], the term \(\int d^3x x^3 \hat{T}_q^{j(\pm)}\) is part of the transverse light-front boost operator and not the transverse light-front rotation operator.

B. Estimates from lattice calculations

In order to determine the quark transverse spin-orbit correlation \[C_s^q\], we need to know four quantities given in Eq. \[39\]. In practice, we can neglect the contribution of the axial FF since it appears multiplied by the mass ratio \(m/3M \approx 10^{-3}\) for \(u\) and \(d\) quarks.

So far, the second Mellin moment of quark transversity GPDs have not yet been extracted from experimental data. We will therefore rely on lattice QCD calculations. In table \[1\] we summarize the results obtained by the QCDSF/UKQCD Collaboration \[31, 32\] for the lowest two Mellin moments of the tensor GPDs \[H_t^q\] and \[E_t^q\]. They are in very good agreement with a more recent calculation by Abdel-Rehim et al. \[33\] which also provides...
TABLE I: Predictions for the scalar charges $\Sigma_q$, tensor charges $\delta_q = \int dx H^q(x, 0, 0)$, anomalous tensor charges $\kappa^q_T = \int dx E^q_T(x, 0, 0)$, and second Mellin moments of $H^q(x, 0, 0)$ and $E^q_T(x, 0, 0)$ for $q = u, d$ from the light-front constituent quark model (LFCQM) and the light-front chiral quark-soliton model (LFQSM) at the scale $\mu^2 \sim 0.26$ GeV$^2$, and from lattice calculations at the scale $\mu^2 = 4$ GeV$^2$.

| Quark model | LFCQM | LFQSM | Lattice | QCDSF/UKQCD Coll. | Abdel-Rehim et al. |
|-------------|-------|-------|---------|----------------|------------------|
| $\delta u$ | 1.165 | 1.241 | 0.857(13) | 0.791(53) |
| $\delta d$ | -0.291 | -0.310 | -0.212(5) | -0.236(33) |
| $\kappa^u_T$ | 3.98 | 3.83 | 2.93(13) | - |
| $\kappa^d_T$ | 2.60 | 2.58 | 1.90(9) | - |
| $\int dx x H^q_T$ | 0.395 | 0.418 | 0.268(6) | 0.264(25) |
| $\int dx x H^q_T$ | -0.099 | -0.105 | -0.052(2) | -0.045(21) |
| $\int dx x E^q_T$ | 1.080 | 1.072 | 0.420(31) | - |
| $\int dx x E^q_T$ | 0.737 | 0.748 | 0.260(23) | - |
| $\Sigma_{u+d}$ | - | - | 8.93(86) | - |
| $\Sigma_{u-d}$ | - | - | 2.20(54) | - |

The numbers in table I can also be used to estimate Burkardt’s correlation. In this case, we obtain from Eq. (66)

$$\langle \delta x J^x_u \rangle_{\text{latt.}} = 0.344, \quad \langle \delta x J^x_d \rangle_{\text{latt.}} = 0.104, \quad (68)$$
$$\langle \delta x J^x_u \rangle_{\text{LFCQM}} = 0.737, \quad \langle \delta x J^x_d \rangle_{\text{LFCQM}} = 0.319, \quad (69)$$
$$\langle \delta x J^x_u \rangle_{\text{LFQSM}} = 0.745, \quad \langle \delta x J^x_d \rangle_{\text{LFQSM}} = 0.321, \quad (70)$$

which can be compared to the values obtained in [37]

$$\langle \delta x J^x_u \rangle_{\text{HYP}} = 0.39, \quad \langle \delta x J^x_d \rangle_{\text{HYP}} = 0.10, \quad (71)$$
$$\langle \delta x J^x_u \rangle_{\text{HO}} = 0.68, \quad \langle \delta x J^x_d \rangle_{\text{HO}} = 0.28, \quad (72)$$

for the hypercentral (HYP) and harmonic oscillator (HO) models.

V. CONCLUSIONS

We introduced and discussed the quark transverse-spin orbit correlation, which is a new piece of information characterizing the nucleon spin structure. We showed that this correlation can be expressed in terms of tensor generalized parton distributions, scalar charges and axial-vector charges. Using results from lattice QCD calculations, we concluded that the quark transverse-spin orbit correlation is very likely negative, just like its longitudinal counterpart. In other words, it is expected that
the quark kinetic orbital angular momentum is in average opposite to the quark spin.

In the process, we compared our quark transverse-spin orbit correlation with Burkardt’s transverse correlation, and obtained several other interesting results. We derived a new sum rule relating twist-2 and twist-3 transversity generalized parton distributions, and also obtained the Wandzura expression for the second Mellin moment of twist-3 transversity generalized parton distributions, which is exact in the chiral limit like in the chiral-even sector. Finally, comparing Ji’s expression for quark kinetic orbital angular momentum to our expression for the quark transverse spin-orbit correlation, we suggested that the scalar charge could be interpreted as a measure of the correlation between quark transversity and transverse spin.

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[38] We would like to correct a remark in [6]: The quark longitudinal spin-orbit operator \( \hat{C}_q^z \) is actually not time-independent since the operator \( \hat{T}^{\mu \nu} \) is in general not conserved. This has however no practical consequences since we are only interested in matrix elements where the initial and final energies are the same.