Four-loop non-singlet splitting functions in the planar limit and beyond

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ABSTRACT: We present the next-to-next-to-next-to-leading order (N^3LO) contributions to the non-singlet splitting functions for both parton distribution and fragmentation functions in perturbative QCD. The exact expressions are derived for the terms contributing in the limit of a large number of colours. For the remaining contributions, approximations are provided that are sufficient for all collider-physics applications. From their threshold limits we derive analytical and high-accuracy numerical results, respectively, for all contributions to the four-loop cusp anomalous dimension for quarks, including the terms proportional to quartic Casimir operators. We briefly illustrate the numerical size of the four-loop corrections, and the remarkable renormalization-scale stability of the N^3LO results, for the evolution of the non-singlet parton distribution and the fragmentation functions. Our results appear to provide a first point of contact of four-loop QCD calculations and the so-called wrapping corrections to anomalous dimensions in \( \mathcal{N} = 4 \) super Yang-Mills theory.

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1 Introduction

Within the gauge theory of the strong interaction, Quantum Chromodynamics (QCD), the precision of theory predictions for hard reactions at colliders crucially depends on our knowledge of hadronic matrix elements for the description of the long-distance hadronic degrees of freedom, once the hard-interaction part due to short-distance physics has been separated by means of QCD factorization. For scattering reactions with initial-state protons the relevant matrix elements are given by the well-known parton distribution functions (PDFs) of the proton, which provide information about the fractions of the proton’s longitudinal momentum carried by the partons.

The dependence of these PDFs on the scale $Q^2$ is generated by evolution equations for the corresponding local operator matrix elements (OMEs). The relevant anomalous dimensions as functions of the Mellin moment $N$, or splitting functions as functions of the momentum fraction $x$, can be computed order by order in perturbative QCD. The corresponding one- and two-loop results have been known since long [1–13]. The current precision is at the three-loop level [14, 15] — see refs. [16–19] for partial recalculations of these results — i.e., at the next-to-next-to-leading order (NNLO), which is nowadays the
accepted standard for analyses of PDFs [20] and forms the backbone of precision predictions at the Large Hadron Collider (LHC).

However, computations for a number key observables at hadron colliders have been performed even at next-to-next-to-next-to-leading order (N^3LO), including the cross section for Higgs-boson production in gluon-gluon fusion [21] and structure functions in deep-inelastic scattering (DIS) [22–25]. The latter results have also found an application in predicting Higgs-boson production in vector-boson fusion at the LHC [26]. Due to QCD factorization, the resulting predictions carry a residual uncertainty and dependence on the factorization scheme due to the missing N^3LO (i.e., four-loop) splitting functions. This situation motivates the computation of the QCD splitting functions at four loops. First steps in this direction have already been taken in refs. [27–31] at low N, and in ref. [32] where large-n_f contributions have been derived at all N.

In the present article, we address the splitting functions for the non-singlet quark evolution equations at four loops in QCD. We use FORCER [33], a FORM [34–36] program for four-loop massless propagators, to compute the anomalous dimensions at fixed integer values of the Mellin variable N. In the planar limit, i.e., for large n_c for a general colour SU(n_c) gauge group, the exact four-loop results for moments up to N = 20 turn out to be sufficient to find and validate the analytic expressions as functions of N in terms of harmonic sums [37, 38] by LLL-based techniques [39, 42] for solving systems of Diophantine equations. Such an approach has been used for anomalous dimensions at the three-and four-loop level before, cf. refs. [32, 43, 44]. Our analytic results in the threshold limit x → 1 (N → ∞) include the (light-like) four-loop cusp anomalous dimension, see ref. [45], which has also been obtained in refs. [46, 47] by different means.

Beyond the large-n_c limit, we have computed the moments up to N = 16 for a general gauge group. These results are insufficient for a reconstruction of the analytic all-N results. They can be used, though, to obtain approximations for the four-loop splitting functions including x-dependent estimates of their residual uncertainties, see, e.g., earlier work at the three-loop level [48–50]. The approximations presented below are sufficiently accurate for the evolution of non-singlet PDFs down to small x, and include numerical results for the non-planar contributions to the four-loop cusp anomalous dimension that are sufficiently precise for phenomenological applications.

For processes with identified hadrons in the final state, QCD factorization requires fragmentation functions (FFs) that account for the physics of hadronization at long distances. Completely analogous to PDFs, the scale dependence of FFs can be computed within perturbative QCD. However, in contrast to the case of initial state hadrons, where the evolution equations for the scale-dependence of the PDFs are controlled by space-like kinematics, \( Q^2 \leq 0 \), the scale evolution of the FFs with \( Q^2 \geq 0 \) requires the so-called time-like splitting functions. These functions are known completely at two loops [9–11, 51–53], see also refs. [54, 55]. The three-loop corrections have been obtained in refs. [56–58] up to a phenomenologically irrelevant small uncertainty in the result for the time-like NNLO quark-gluon splitting function. First NNLO analyses of FFs have been performed recently [59, 60].

\[ 1 \text{See also ref. [40], summarized in [41], pg. 16.} \]
The three-loop results in refs. [56–58] have been derived using well-known relations between space- and time-like kinematics, i.e., the Drell-Yan-Levy relation for the analytic continuation in energy $q^2 \rightarrow -q^2$ and the Gribov-Lipatov relation in $x$-space [61, 62], see also refs. [63, 64], and generalizations based on conformal symmetry yielding a universal reciprocity-respecting evolution kernel [65–67]. Exploiting these relations, it is possible to use (space-like) DIS results to predict (time-like) cross sections for single-particle inclusive electron-positron annihilation. Thus, we are able to present here also the flavour non-singlet evolution equations for FFs at four loops in QCD.

This article is organized as follows. In section 2 we specify our notations and present the theoretical framework for obtaining our results. In particular we address the basis of non-singlet operators, their renormalization and the respective anomalous dimensions. We sketch the work-flow of the perturbative computation up to four loops, list all colour factors to this order and discuss general and end-point properties of the anomalous dimensions and splitting functions.

In section 3 we present the results of our fixed-$N$ diagram calculations of the four-loop non-singlet anomalous dimensions and their all-$N$ generalization in the large-$n_c$ limit. We discuss the large-$N$ behaviour of the latter which includes the four-loop cusp anomalous dimension. The $x$-space counterparts of these anomalous dimensions, i.e., the splitting functions, are addressed in section 4. We present the exact formulae and compact parametrizations for the large-$n_c$ splitting functions, and approximate expressions for all cases that cannot be obtained exactly for now.

Two important applications of these results are presented in section 5: we present high-accuracy numerical results for large-$x$ coefficients, in particular the four-loop cusp anomalous dimension in QCD, and illustrate the $N^3$LO evolution of all three types of non-singlet quark distributions. The $N^3$LO non-singlet evolution is extended to the ‘time-like’ case of final-state fragmentation functions in section 6. We summarize our main results and provide a brief outlook in section 7.

The appendices contain the Feynman rules in appendix A, the exact results for the anomalous dimensions at $1 \leq N \leq 16$ at four loops in appendix B, and the analytic expression for the difference of the time-like and space-like four-loop splitting functions in appendix C. Finally appendix D provides the complete all-$N$ result for the terms with $\zeta_5$, which may be of theoretical interest.

## 2 Theoretical framework and calculations

The standard set of spin-$N$ twist-two irreducible flavour non-singlet quark operators is given by

$$O^{\text{ns}}_{\{\mu_1, \ldots, \mu_N\}} = \overline{\psi} \lambda^\alpha \gamma_{\{\mu_1} D_{\mu_2} \ldots D_{\mu_N\}} \psi, \quad \alpha = 3, 8, \ldots, (n_f^2 - 1), \quad (2.1)$$

where $\psi$ represents the quark field, $D_\mu = \partial_\mu - igA_\mu$ the covariant derivative, and $\lambda^\alpha$ the diagonal generators of the flavour group SU($n_f$). It is understood in eq. (2.1) that the symmetric and traceless part is taken with respect to the Lorentz indices $\mu_i$ in the curly brackets.
Figure 1. Vertices with additional gluons arising from the operators $O_{\{\mu_1,...,\mu_N\}}^{ns}$ in perturbative QCD. Vertices with up to $L$ gluons need to be considered at $L$ loops at Mellin moments $N > L$.

We consider (spin-averaged) matrix elements of these operators (OMEs), specifically

$$\langle p_1 | O_{\{\mu_1,...,\mu_N\}}^{ns} | p_2 \rangle ,$$

for external quark (or anti-quark) fields with momenta $p_1$ and $p_2$. The operators $O^{ns}$ in eq. (2.2) are contracted with tensors of rank $N$,

$$\Delta^{\mu_1} \ldots \Delta^{\mu_N} ,$$

where $\Delta$ is a light-like vector, $\Delta^2 = 0$. In the present context we need to compute OMEs of renormalized operators with zero momentum flow through the operator vertex, thus $p_1 = p_2 = p$ in eq. (2.2) for the (off-shell, $p^2 \neq 0$) momenta of the external (anti-)quarks,

$$[A^{ns}](N) = \Delta^{\mu_1} \ldots \Delta^{\mu_N} \langle p | [O^{ns}]_{\{\mu_1,...,\mu_N\}} | p \rangle .$$

Here and below we use square brackets $[\ldots]$ to denote renormalized operators (in a minimal subtraction scheme $[68, 69]$ of dimensional regularization $[70, 71]$). This reduces the vertex diagrams for the OMEs to quark two-point functions and, therefore, the computational complexity to propagator-type diagrams. The perturbative expansion of the operator in eq. (2.1) contracted with eq. (2.3) generates vertices with additional gluons as depicted in figure 1. The current four-loop calculation requires up to four additional gluons. The corresponding Feynman rules are presented in appendix A, see refs. $[5, 72]$ for earlier calculations at two- and three-loop accuracy.

In order to derive the anomalous dimensions for the scale dependence of the non-singlet PDFs we need to perform the renormalization of those operators $O^{ns}$, which proceeds multiplicatively as

$$[O^{ns}] = Z_{ns} O^{ns} .$$

The anomalous dimensions $\gamma_{ns}$ governing the scale dependence of these operators,

$$\frac{d}{d \ln \mu^2} [O^{ns}] = - \gamma_{ns} [O^{ns}]$$

are connected to the factors $Z_{ns}$ in eq. (2.5) by

$$\gamma_{ns} = - \left( \frac{d}{d \ln \mu^2} Z_{ns} \right) Z_{ns}^{-1} .$$

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All flavour differences of quark-anti-quark sums (+) and differences (−) evolve in with the same anomalous dimensions \( \gamma_{ns}^+(N) \) and \( \gamma_{ns}^-(N) \), and the total valence distribution with \( \gamma_{ns}^\nu(N) = \gamma_{ns}^+(N) + \gamma_{ns}^-(N) \), see, e.g., ref. [14]. These quantities are related to the corresponding splitting functions \( P_{ns}^\pm(x) \) and \( P_{ns}^a(x) \) by a Mellin transform,

\[
\gamma_{ns}(N) = -\int_0^1 dx \, x^{N-1} P_{ns}(x) ,
\]

where the relative sign is a standard convention. In perturbation theory these quantities can be expanded in powers of the strong coupling constant \( \alpha_s \). Here and below we normalize \( a_s \equiv \alpha_s/(4\pi) \), so that up to four loops

\[
\gamma_{ns}(N) = a_s \gamma_{ns}^{(0)}(N) + a_s^2 \gamma_{ns}^{(1)}(N) + a_s^3 \gamma_{ns}^{(2)}(N) + a_s^4 \gamma_{ns}^{(3)}(N) ,
\]

and similarly for the splitting functions \( P_{ns}(x) \) and other quantities. The first-order quantity \( \gamma_{ns}^{(0)} \) is the same for all three cases given above. \( \gamma_{ns}^+ \) and \( \gamma_{ns}^- \) differ at order \( a_s^2 \), and a non-vanishing flavour-independent (‘sea’) contribution \( \gamma_{ns}^s \) occurs at order \( a_s^3 \) for the first time [14]. The fourth-order contributions \( \gamma_{ns}^{(3)}(N) \) to all three quantities are addressed in the present article.

The actual computation follows a well-established production chain. The Feynman diagrams for the OMEs in eq. (2.2) are generated up to four loops using QGRAF [73]. The latest version [74] of the symbolic manipulation program FORM [34, 35] and its multi-threaded version TFORM [36] are used for all further steps. The QGRAF output is processed by a program that assigns the topology and computes the colour factor using the code of ref. [75]; the group invariants occurring in the present case are listed in table 1. Diagrams of the same topology and colour factor are combined to meta diagrams for computational efficiency, where lower-order self-energy insertions are treated as described in ref. [76]. Considering all color factors, this procedure leads to 1 one-loop, 7 two-loop, 53 three-loop and 650 four-loop meta diagrams for \( \gamma_{ns}^\pm \); and 1 three-loop and 29 four-loop meta diagrams for \( \gamma_{ns}^s \). For comparison: the output of QGRAF consists of 15901 four-loop diagrams. The running of the meta diagrams is managed using the database program MINOS [77].

The diagram calculations are done in dimensional regularization [70, 71] with the FORCER program [33] which was already used for the \( N \leq 6 \) and high-\( n_f \) computations in refs. [31, 32]. Our agreement (after renormalization, see below) with those results, which were obtained in a different theoretical framework, provides a strong check of our present setup. The FORCER program itself has been validated in calculations of the four-loop renormalization of Yang-Mills theories to all powers of the gauge parameter, see ref. [79], and has recently been applied — together with the algorithms for the \( R^s \) operation [80, 81] developed in ref. [82] — in five-loop computations of the beta function, Higgs-boson decays to hadrons and the \( R \)-ratio in \( e^+e^- \)-annihilation in refs. [83, 84].

The bare results \( A_{ns} \) for the OMEs in eq. (2.4) obtained in this way are then subject to renormalization which we perform in the standard modified minimal subtraction scheme \( \overline{\text{MS}} \) [68, 69]. In this scheme and in \( D = 4 - 2\epsilon \) dimensions the strong coupling \( \alpha_s \) evolves according to

\[
\frac{d}{d\ln \mu^2} \frac{\alpha_s}{4\pi} = \frac{d a_s}{d \ln \mu^2} = \beta(a_s) = -\epsilon a_s - \beta_0 a_s^2 - \beta_1 a_s^3 - \beta_2 a_s^4 - \ldots ,
\]

(2.10)
Table 1. The colour factors for non-singlet OMEs up to four loops with their numerical values in SU($n_c$) with $N_R = n_c$ and QCD, see also ref. [78] for a discussion on the normalization of $d^{abc}d_{abc}$. The second column gives the notations of the result files, which are distributed with this article on https://arxiv.org. As in many other articles, we suppress the colour factor $T_F (~= 1/2$ in SU($n_c$)) which can be readily re-instated.

where $\beta(a_s)$ denotes the usual four-dimensional beta function in QCD, with coefficients $\beta_0 = 11/3 C_A - 2/3 n_f$ etc, and $n_f$ represents the number of active quark flavours.

Using eq. (2.5) the renormalized OMEs $[A^{ns}]$ are obtained by

$$[A^{ns}](N) = Z_\psi Z_{ns}(N) A^{ns}(N),$$

where we have made all dependences on $N$ explicit. The factor $Z_\psi$ denotes the quark wave function renormalization constant accounting for the external quarks field with off-shell momenta in eq. (2.4), see, for instance ref. [79]. Unlike $Z_{ns}$, the quantities $Z_\psi$ and $A^{ns}$ are gauge-dependent, hence also the renormalization constant $Z_\xi$ of the gauge parameter is required.

The resulting operator renormalization factors $Z_{ns}$ in eq. (2.5) can be expressed as a Laurent series in $\varepsilon$ as

$$Z_{ns} = 1 + a_s \left\{ \frac{1}{\varepsilon} \gamma_{ns}^{(0)} \right\} + a_s^2 \left\{ \frac{1}{\varepsilon^2} \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \gamma_{ns}^{(0)} \right)^2 - \frac{1}{2} \gamma_{ns}^{(0)} \right) \right) \right) \right) \right\} + \frac{1}{2\varepsilon} \gamma_{ns}^{(1)} \right\}$$

$$+ a_s^3 \left\{ \frac{1}{\varepsilon^3} \left( \frac{1}{6} \gamma_{ns}^{(0)} \right)^3 - \frac{1}{2} \beta_0 \left( \frac{1}{2} \gamma_{ns}^{(0)} \right)^2 + \frac{1}{3} \beta_0^2 \gamma_{ns}^{(0)} \right\} + \frac{1}{2\varepsilon} \gamma_{ns}^{(1)} \right\}$$

$$+ \frac{1}{3\varepsilon} \gamma_{ns}^{(2)} \right\} + a_s^4 \left\{ \frac{1}{\varepsilon^4} \left( \frac{1}{12} \gamma_{ns}^{(0)} \right)^4 - \frac{1}{4} \beta_0 \left( \gamma_{ns}^{(0)} \right)^3 + \frac{11}{24} \beta_0^2 \left( \gamma_{ns}^{(0)} \right)^2 - \frac{1}{4} \beta_0^3 \gamma_{ns}^{(0)} \right\}$$

$$+ \frac{1}{3\varepsilon} \gamma_{ns}^{(1)} \right\} + a_s^5 \left\{ \frac{1}{\varepsilon^5} \left( \frac{1}{24} \gamma_{ns}^{(0)} \right)^5 - \frac{7}{12} \beta_0 \left( \gamma_{ns}^{(0)} \right)^4 + \frac{1}{4} \beta_0^2 \gamma_{ns}^{(0)} + \frac{1}{2} \beta_0 \beta_1 \gamma_{ns} + \frac{1}{3} \beta_1 \left( \gamma_{ns}^{(0)} \right)^2 \right\}$$

$$+ \frac{1}{3\varepsilon} \gamma_{ns}^{(1)} \right\} + \frac{1}{8} \left( \gamma_{ns}^{(0)} \right)^2 - \frac{1}{4} \gamma_{ns}^{(2)} \beta_0 - \frac{1}{4} \beta_0 \gamma_{ns}^{(1)} - \frac{1}{4} \beta_2 \gamma_{ns}^{(0)} \right\} + \frac{1}{4\varepsilon} \gamma_{ns}^{(3)} \right\}. \quad (2.12)$$

In this manner, the anomalous dimensions $\gamma_{ns}$ have been computed for a general gauge group at $1 \leq N \leq 16$, i.e, $\gamma_{ns}^+$ at even $N$ and $\gamma_{ns}^-$ at odd $N$. The exact results are listed
in appendix B; numerical values for QCD can be found in section 3. The hardest (non-planar) diagrams do not contribute in the limit of a large number of colours $n_c$, where the functions $\gamma_{n_s}^+(N)$ and $\gamma_{n_s}^-(N)$ are identical, as it is evident from diagrammatical analyses and the known $x$-space expressions for $P_{ns}^\pm(x)$, see refs. [14, 32, 66, 85]. Consequently we were able to obtain the even-$N$ and odd-$N$ values of the large-$n_c$ anomalous dimension, which is structurally simpler than full QCD results, even up to $N = 20$.

So far, fixed-$N$ values of anomalous dimensions have been found to be fractions of (large) integer numbers, multiplied at most by values $\zeta_3 \ldots \zeta_{2L-3}$ of the Riemann zeta-function at $L$ loops. The denominator structure of the fractions suggests analytic all-$N$ expressions in terms of harmonic sums [37, 38] up to weight $2L - 1$. Assuming no numerator-$N$ terms, cf. refs. [22–24], the most complicated parts (without a factor $\zeta_n$) of the non-singlet anomalous dimensions at $n$ loops read

$$\gamma_{n_s}^{(n)}(N) = \sum_{w=0}^{2n+1} c_{00w} S_w(N) + \sum_a \sum_{k=1}^{2n+1} \sum_{w=0}^{2n+1-k} c_{akw} D_a^k S_w(N) \ ,$$

(2.13)

where $D_a^k$ are simple denominators,

$$D_a^k = (N + a)^{-k} \ ,$$

(2.14)

and $S_w(N)$ is a shorthand for all harmonic sums of a given weight $w$ with $S_0(N) \equiv 1$. The calculated moments suggest $a = 0, 1$ for $\gamma_{n_s}^{(3)}(N)$, as at three loops [14] and for the $n_f^2$ and $n_f^3$ four-loop contributions [32]. The function $\gamma_{n_s}^8(N)$, on the other hand, includes terms with $a = -1$ and $a = 2$.

The functions $\gamma_{n_s}^{(3)}(N)$ contain harmonic sums up to weight $w = 7$, hence the ansatz (2.13) includes far too many unknown coefficients $c_{akw}$ for a direct determination from the (small) number of calculated moments. However, these coefficients are integer modulo some predictable powers of 2 and 3. Therefore the systems of equations derived from eq. (2.13) can be turned into Diophantine systems which require far fewer equations than unknowns and which can be solved by LLL-based techniques [39–42]. This approach has been successfully applied before in refs. [32, 43, 44].

In this context it is crucial to constrain eq. (2.13) as far as possible based on general properties of the anomalous dimensions $\gamma_{n_s}(N)$. Here three issues are worth pointing out. First, the functional forms of the $\gamma_{n_s}(N)$ are (conjectured to be) constrained by ‘self-tuning’ [66, 67],

$$\gamma_{n_s}(N) = \gamma_u (N + \sigma \gamma_{n_s}(N) - \beta(a_s/a_u)) \ ,$$

(2.15)

where $\sigma = -1(1)$ for the space-like (time-like) anomalous dimensions, and the non-singlet universal evolution kernel $\gamma_u$ is reciprocity-respecting (RR), i.e., invariant under the replacement $N \rightarrow (1 - N)$. By expanding the r.h.s. of eq. (2.15) about $N$ and inserting the perturbation series of all quantities involved, $\gamma_u$ can be expressed in terms of the $\overline{\text{MS}}$ anomalous dimensions, see also ref. [86]. Expressing the latter in terms of $\gamma_0 = \gamma_{n_s}^{S(n)} = \gamma_{n_s}^{T(n)}$ and the average of the space-like and time-like expansion coefficients $\overline{\gamma}_n =
where we have used the abbreviation \( d_N = d/dN \) and suppressed the \( N \)-dependencies for brevity. A convenient way to take these derivatives is via inverse Mellin transforms to \( x \)-space, where the multiplication with \( \ln^n x \) corresponds to the \( N \)-space operator \( d^n/dN^n \), and Mellin transforms of the result. The required manipulations can be readily performed using algorithms for harmonic sums, harmonic polylogarithms and their (inverse) Mellin transformations \[37, 87, 88\] which have been implemented in publicly available FORM packages described in ref. \[34\].

Since the difference between the time-like and space-like anomalous dimensions is known to four loops, eq. (2.13) can be applied to the RR quantity \( \gamma^{(3)}_u \) instead of \( \gamma^{(3)}_{ns} \). This implies that the denominators \( 1/N \) and \( 1/(N+1) \) can only enter in the combination \( 1/(N(N+1)) \), and that only RR (combinations of) harmonic sums occur, see refs. \[89, 90\], which reduces the number of sums at weight \( w \) from \( 2 \cdot 3^{w-1} \) to \( 2^{w-1} \). Assuming that only powers of \( 1/(N+1) \) enter in addition, the total number of basis functions in eq. (2.13) up to weight \( w \) is \( 2^{w-1} - 1 \), e.g., 255 for \( w = 7 \). Even taking account end-point constraints, see below, this is a prohibitively large number for now.

Second, the identical leading-\( n_c \) terms of \( \gamma^{(3)}_{ns}(N) \) contain only non-alternating harmonic sums, i.e., only positive indices in eq. (3.4). This reduces the number of RR sums of weight \( w \) to the Fibonacci number \( F(w) \), i.e., 1, 1, 2, 3, 5, 8, 13 for \( w = 1 \) to \( w = 7 \), as can be seen by counting the number of binomial harmonic sums at weight \( w \) \[89\]. Considering all combinations with additional powers of the weight-1 object \( 1/(N(N+1)) \), the total number of functions up to weight \( w \) in eq. (2.13) amounts to \( F(w+4) - 2 \), e.g., 87 for \( w = 7 \).

The third and final point is that the \( N \to \infty \) (large-\( x \)) and \( N \to 0 \) (small-\( x \)) limits of the anomalous dimensions (splitting functions) provide a substantial number of constraints. If one disregards terms of order \( O(1/N^2) \) for \( N \to \infty \), then all three non-singlet anomalous dimensions \( \gamma^{(a)}_{ns}(N) \), \( a = \pm, - \), \( v \), are identical and given by \[65\] \( \gamma_v \) is the Euler-Mascheroni constant

\[
\gamma^{(n-1)}_{ns}(N) = A_n (\ln N + \gamma_v) + B_n + C_n (\ln N + \gamma_v) N^{-1} - \left( \tilde{D}_n - A_{n+1} \right) N^{-1} .
\]  

Here the coefficients \( A_n \) — the \( n \)-loop (light-like) cusp anomalous dimension — and \( B_n \) provide genuine \( n \)-loop information. The coefficients \( C_n \) and \( \tilde{D}_n \), on the other hand, can be expressed in terms of lower-order information (see eqs. (3.10) and (3.11) below). This and the absence of second and higher powers of \( \ln N \) in eq. (2.17), and similar if less stringent constraints on \( N^{-k} \ln N \) terms with \( k > 1 \), provide a substantial number of constraints on the coefficients in eq. (2.13).
The small-$x$ expansion of the splitting function $P_{ns}^{(n)}(x)$ shows a double-logarithmic enhancement, i.e., there are contributions of the form $x^a \ln^b x$ with $a > 0$ and $b \leq 2n$. The leading-logarithmic (LL) contributions to $P_{ns}^\pm$ have been known to all orders for a long time [91, 92]. This resummation has been extended to next-to-next-to-logarithmic accuracy for the $x^{2k} \ln^b x$ contributions to $P_{ns}^+$ and the $x^{2k+1} \ln^b x$ contributions to $P_{ns}^-$ at all $k \geq 0$ [93, 94]. The formal structure of these results is analogous to their time-like counterparts [95, 96], but the numerical pattern is completely different such that the space-like resummation is of no direct phenomenological use. The functions $P_{ns}^+(x)$ and $P_{ns}^-(x)$ are the same in the large-$n_c$ limit, hence in this case the small-$x$ resummation constrains the coefficients contributing to

$$x^a \ln^b x \quad \text{for} \quad a \geq 0 \quad \text{and} \quad 4 \leq b \leq 6 \quad . \quad (2.18)$$

An alternative approach to the limit $N \to 0$, i.e., the small-$x$ logarithms for $a = 0$, has been pursued in ref. [97]. In the large-$n_c$ limit, the generalization

$$\gamma_{ns}(N, a_s) \cdot (\gamma_{ns}(N, a_s) + N - \beta(a_s)/a_s) = O(1) \quad (2.19)$$

of the LL relation in refs. [91, 92] correctly (re-) produces all $a = 0$ small-$x$ logarithms obtained in refs. [14, 32, 93], after correcting typos in eqs. (25) and (26) of ref. [97]. Hence we can assume that eq. (2.19) is also correct for the $n_f^0$ and $n_f^1$ four-loop contributions at large $n_c$.

Together these relations comprise 18 large-$N$ and 28 small-$x$ constraints for the $n_f^0$-terms at four loops eliminating more than half of the 87 free parameters of the $w=7$ large-$n_c$ ansatz, after which it is possible to solve the remaining system of Diophantine equation using the moments $N = 1, \ldots, 18$ with the program axb() of the CALC package [42]. The resulting analytic expressions for $\gamma_{ns}^{(3)}(N)$ agree with the result of the diagram calculations at $N = 19$ and $N = 20$. This agreement renders it extremely likely — although, of course, not mathematically certain — that these results (and, therefore, the above structural conjectures and features used in their derivation) are correct.

As mentioned above, present information and understanding appears not to be sufficient for extending these analytic results beyond the large-$n_c$ contributions for the $n_f^0$ and $n_f^1$ parts of $\gamma_{ns}^{(3)a}(N)$ for any $a = +, -, s$. For the remaining functions we resort to $x$-space approximations based on the first eight even-$N$ or odd-$N$ moments supplemented by the large-$x$ and small-$x$ constraints discussed above. These approximations and their error estimates can be constructed in the same manner as those for the three-loop splitting functions in refs. [48–50]. The present results are more accurate, though, due to the higher number of available moments and the improved understanding of the end-point limits. The fact that the large-$N$ limit (2.17) includes only the two free parameters $A_4$ and $B_4$, in particular, results in a high accuracy of these coefficients which are relevant also in the context of the soft-gluon exponentiation, see refs. [98–101] and references therein, and beyond.
3 Results in $N$-space

We start presenting our results by writing down the moments to $N = 16$ of the non-singlet four-loop anomalous dimensions for QCD in a numerical form. The exact results for a general gauge group with one set of fermions can be found in appendix B. For $\gamma_{ns}^{3+}(N)$ we separately display the leading (subscript $L$) and non-leading (subscript $N$) contributions in the large-$n_c$ limit of SU($n_c$) at $n_c = 3$. The former correspond to the colour factors $C_F n_c^{3-k} n_f^k$. The latter collect all other terms, which are suppressed by two or more powers of $n_c$, cf. table 1 above.

The first eight even-$N$ values of $\gamma_{ns}^{3+}(N)$, normalized as in eq. (2.9) above — division by $2.5 \cdot 10^4$ provides an approximate conversion to an expansion in $\alpha_s$ — are given by

$$\gamma_{ns}^{3+}(2) = 15079.979904_L + 1583.245584_N - \left(3610.4544273_L + 137.0956130_N\right) n_f + (114.04980426_L + 3.73169344_N) n_f^2 + 1.908433871_L n_f^3,$$

$$\gamma_{ns}^{3+}(4) = 26818.645638_L + 996.7898428_N - \left(6549.4629888_L + 154.7093878_N\right) n_f + (231.3355779_L + 7.35324122_N) n_f^2 + 3.884444957_L n_f^3,$$

$$\gamma_{ns}^{3+}(6) = 33391.136191_L + 801.6538203_N - \left(8219.2929199_L + 167.0140214_N\right) n_f + (299.1252425_L + 9.35853808_N) n_f^2 + 5.073482406_L n_f^3,$$

$$\gamma_{ns}^{3+}(8) = 38122.913329_L + 685.2274355_N - \left(9424.0745879_L + 175.4883891_N\right) n_f + (34.75255308_L + 10.77524037_N) n_f^2 + 5.937491139_L n_f^3,$$

$$\gamma_{ns}^{3+}(10) = 41872.765979_L + 600.8875674_N - \left(10377.826722_L + 181.7070215_N\right) n_f + (385.45469324_L + 11.87446948_N) n_f^2 + 6.618723172_L n_f^3,$$

$$\gamma_{ns}^{3+}(12) = 44999.013143_L + 533.8372124_N - \left(1171.437026_L + 186.4798090_N\right) n_f + (416.75989357_L + 12.77304489_N) n_f^2 + 7.181732766_L n_f^3,$$

$$\gamma_{ns}^{3+}(14) = 47688.997330_L + 477.6377749_N - \left(11852.885145_L + 190.2641681_N\right) n_f + (443.46765303_L + 13.53286843_N) n_f^2 + 7.661731916_L n_f^3,$$

$$\gamma_{ns}^{3+}(16) = 50054.446557_L + 428.9331378_N - \left(12450.928224_L + 193.3399276_N\right) n_f + (466.7845732_L + 14.19090917_N) n_f^2 + 8.080152509_L n_f^3,$$

and the first eight odd-$N$ values of $\gamma_{ns}^{3-}(N)$ read

$$\gamma_{ns}^{3-}(1) = 0,$$

$$\gamma_{ns}^{3-}(3) = 22101.772707_L + 999.394009_N - \left(5359.5335080_L + 139.7317364_N\right) n_f + (183.0229066_L + 5.91723185_N) n_f^2 + 3.058631170_L n_f^3,$$

$$\gamma_{ns}^{3-}(5) = 30431.615620_L + 841.7002595_N - \left(7466.3287550_L + 160.0856189_N\right) n_f + (268.62851396_L + 8.47774637_N) n_f^2 + 4.534727508_L n_f^3,$$

$$\gamma_{ns}^{3-}(7) = 35914.342449_L + 722.2583518_N - \left(8861.7720221_L + 171.1596475_N\right) n_f + (325.00219825_L + 10.13425314_N) n_f^2 + 5.534366480_L n_f^3,$$

$$\gamma_{ns}^{3-}(9) = 40092.293568_L + 632.8933507_N - \left(9925.1782168_L + 178.6640945_N\right) n_f + (367.49639662_L + 11.36914948_N) n_f^2 + 6.295934539_L n_f^3,$$

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\begin{align}
\gamma^{(3)}_{\text{ns}}(11) &= 43499.696829_L + 561.6966621_N - (10791.034298_L + 184.1995847_N) n_f \\
&+ (401.78232680_L + 12.35578151_N) n_f^2 + 6.912368819_L n_f^3 , \\
\gamma^{(3)}_{\text{ns}}(13) &= 46390.354670_L + 502.2596221_N - (11524.077148_L + 188.4837700_N) n_f \\
&+ (430.5996163_L + 13.17711657_N) n_f^2 + 7.430544877_L n_f^3 , \\
\gamma^{(3)}_{\text{ns}}(15) &= 48907.084426_L + 451.0029039_N - (12160.990233_L + 191.9086558_N) n_f \\
&+ (455.49394328_L + 13.88095204_N) n_f^2 + 7.877634036_L n_f^3 . \tag{3.2}
\end{align}

It is clear from these results, that the large-\(n_c\) limit alone provides an excellent approximation to the individual \(n_f^2\) coefficients except for the lowest values of \(N\). The non-large-\(n_c\) ‘correction’ amounts to 10% and 4% for the \(n_f^0\) and \(n_f^1\) terms, respectively, but 2% or less at \(N \geq 7\) in both cases.

We have computed the first nine odd-\(N\) values of the ‘sea’ contribution \(\gamma^{(3)s}_{\text{ns}}(N)\) to the four-loop anomalous dimension for the overall valence distribution, and find

\begin{align}
\gamma^{(3)s}_{\text{ns}}(1) &= 0 , \\
\gamma^{(3)s}_{\text{ns}}(3) &= 18.9700898832 n_f - 1.6109396433 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(5) &= 9.1402406178 n_f + 0.1525610933 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(7) &= 4.0470106556 n_f + 0.3095914493 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(9) &= 1.9658456985 n_f + 0.2887522942 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(11) &= 1.0102117327 n_f + 0.2474854100 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(13) &= 0.5255291656 n_f + 0.2099342070 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(15) &= 0.2612077170 n_f + 0.1791623892 n_f^2 , \\
\gamma^{(3)s}_{\text{ns}}(17) &= 0.1094470531 n_f + 0.1544243611 n_f^2 . \tag{3.3}
\end{align}

We now turn to the analytic all-\(N\) expressions for the \(n_f^0\) and \(n_f^1\) parts of the four-loop non-singlet anomalous dimensions \(\gamma^{(3)\pm}_{\text{ns}}\) in the large-\(n_c\) limit. The complete lower-order contributions can be found, in a different notation but the same normalization, in eqs. (3.4)–(3.8) of ref. [14]. The anomalous dimensions can be expressed in terms of the denominators \(D_{\eta}^{(k)}\) in eq. (2.14) and harmonic sums [37, 38] at argument \(N\), which are recursively defined by

\begin{align}
S_{\pm m}(N) &= \sum_{n=1}^{N} (\pm 1)^n n^{-m} , \\
S_{\pm m_1, m_2, ..., m_d}(N) &= \sum_{n=1}^{N} (\pm 1)^n n^{-m_1} S_{m_2, ..., m_d}(n) . \tag{3.4}
\end{align}

The weight \(w\) of the harmonic sums is defined by the sum of the absolute values of the indices \(m_d\). Sums up to \(w = 2n - 1\) occur in the \(n\)-loop anomalous dimensions. The argument \(N\) of the sums is suppressed for brevity below, and we use the shorthand \(\eta = 1/(N(N+1)) = D_0 D_1\).
The identical large-$n_c$ parts of the functions $\gamma_{\text{ns}}^{(3)+}(N)$ and $\gamma_{\text{ns}}^{(3)-}(N)$ are given by

$$
\gamma_{\text{ns},L}(N) = C_F \left( n_c^2 \gamma_{L,1}^{(3)}(N) + n_c \eta_f^2 \gamma_{L,1}^{(3)}(N) + n_c \eta_f^2 \gamma_{L,2}^{(3)}(N) + n_f^2 \gamma_{L,3}^{(3)}(N) \right), \tag{3.5}
$$

where the $n_f^2$ and $n_f^2$ contributions to eq. (3.5) have been given in eqs. (3.1) and (3.6) of ref. [32]; the latter has first been derived in ref. [102]. Our new results are

$$
\gamma_{L,0}^{(3)}(N) =
\begin{align*}
&16 \left( 1379569/82944 - 7453/24 D_1^2 - 102641/648 D_1^2 - 9883/864 D_1^2 - 209/12 D_1^5 \\
&-95/12 D_1^5 - 5/4 D_1^3 + 534767/5184 \eta - 231341/2592 \eta^2 - 15469/1296 \eta^3 - 12997/864 \eta^4 \\
&-83/12 \eta^5 - 25/24 \eta^6 - 5/8 \eta^7 - 55/4 \zeta_5 + 40 \zeta_5 \eta - 25 \zeta_5 \eta^2 - 1517/144 \zeta_3 + 839/108 \zeta_3 \eta \\
&+13/24 \zeta_3 \eta^2 - 2/3 \zeta_3 \eta^3 - 1/2 \zeta_3 \eta^4 + 42139/648 S_1 - 130795/648 S_1 D_1^2 - 298/9 S_1 D_1^3 \\
&-995/24 S_1 D_1^4 - 92/3 S_1 D_1^5 - 15/2 S_1 D_1^6 + 278627/1296 S_1 \eta + 19757/2592 \eta^2 \\
&+3625/48 S_1 \eta^3 + 1789/72 S_1 \eta^4 + 9 S_1 \eta^5 + 7/2 S_1 \eta^6 + 422/27 S_1 \zeta_3 - 62/3 S_1 \zeta_3 \eta \\
&-7/3 S_1 \zeta_3 \eta^2 + 105/3 S_1 \zeta_3 \eta^3 - 24211/432 S_2 + 23153/216 S_2 \eta + 143/3 S_2 \eta^2 + 165 S_2 D_1^2 \\
&+107 S_2 D_1^2 - 18725/1296 S_2 \eta + 23689/432 S_2 \eta^2 + 3187/144 S_2 \eta^3 + 229/12 S_2 \eta^4 + 29/4 S_2 \eta^5 \\
&-20 S_2 \eta^6 + 8 S_3 \zeta_3 \eta^2 + 17591/864 S_3 - 5099/72 S_3 D_1^2 - 5/3 S_3 D_1^3 - 2 S_3 D_1^4 + 4373/108 S_3 \eta \\
&+1051/72 S_3 \eta^2 + 40/3 S_3 \eta^3 + 27/4 S_3 \eta^4 - 13/3 S_3 \zeta_3 + 2 S_3 \zeta_3 \eta - 55291/864 S_4 \\
&+155/6 S_4 D_1^2 + 7 S_4 D_1^3 + 3233/144 S_4 \eta + 103/12 S_4 \eta^2 + 11/2 S_4 \eta^3 + 4 S_4 \zeta_3 + 227/4 S_5 \\
&-12 S_5 D_1^2 + 1/3 S_5 \eta + 7 S_5 \eta^2 - 65/3 S_6 + 10 S_6 \eta + 10 S_6 \eta^2 - 56/3 S_6 D_1^2 - 71/3 S_6 D_1^3 \\
&-35 S_6 D_1^4 - 20 S_6 \eta + 23/12 S_6 \eta^2 - 587/8 S_6 \eta^3 - 115/4 S_6 \eta^4 - 109/3 S_6 \eta^5 \\
&-11 S_6 \eta^6 + 40 S_6 \zeta_3 \eta^2 - 165 S_6 \zeta_3 \eta^3 + 5529/81 S_7 + 184/3 S_7 \eta^2 + 2 S_7 \eta^3 + 4 S_7 \zeta_3 + 227/4 S_5 \\
&-38 S_7 \eta^4 + 12 S_7 \eta^5 - 241/12 S_7 \eta^6 - 137/6 S_7 \eta^7 - 17/2 S_7 \eta^8 + 890/27 S_7 \eta^9 - 22/3 S_7 \eta^{10} \\
&+53/3 S_7 \eta^{11} - 79/6 S_7 \eta^{12} - 115 S_7 \eta^{13} - 4 S_7 \zeta_3 \eta - 4745/72 S_7 \eta + 143 S_7 \eta^2 D_1^2 - 46/3 S_7 \eta^3 \\
&-S_7 \eta^4 - 70/3 S_7 \eta^5 - 4 S_7 \eta^6 - 20 S_7 \eta^7 - 5529/81 S_7 + 184/3 S_7 \eta^2 + 2 S_7 \eta^3 + 4 S_7 \zeta_3 + 227/4 S_5 \\
&-38 S_7 \eta^4 + 12 S_7 \eta^5 - 241/12 S_7 \eta^6 - 137/6 S_7 \eta^7 - 17/2 S_7 \eta^8 + 9029/216 S_7 \eta^9 - 19 S_7 \eta^{10} \\
&-65 S_7 \eta^{11} - 38 S_7 \eta^{12} - 183 S_7 \eta^{13} - 9 S_7 \eta^{14} - 3635/72 S_7 \eta + 12 S_7 \eta^2 D_1^2 + 5/3 S_7 \eta^3 \\
&-10 S_7 \eta^4 + 31/2 S_7 + 3 S_7 \eta^2 - 16 S_7 + 925/18 S_7 + 76/3 S_7 \eta^2 - 8 S_7 \eta^3 \\
&-1675/36 S_7 \eta^4 - 21 S_7 \eta^5 - 4 S_7 \eta^6 - 4 S_7 \zeta_3 \eta - 5581/72 S_7 + 22 S_7 \eta^3 D_1^2 - 53/3 S_7 \eta^4 \\
&-16 S_7 \eta^5 + 143/6 S_7 \eta^6 - 11 S_7 \eta^7 + 14 S_7 \eta^8 - 6899/72 S_7 + 24 S_7 \eta^3 D_1^2 - 74/3 S_7 \eta^4 \\
&-11 S_7 \eta^5 + 57/2 S_7 \eta^6 - 25 S_7 \eta^7 - 26 S_7 \eta^8 + 6 S_7 \eta^9 - 36 S_7 \eta^{10} - 28 S_7 \eta^{11} \\
&-12 S_7 \eta^{12} + 12 S_7 \eta^{13} + 24 S_7 \eta^{14} + 6 S_7 \eta^{15} + 18 S_7 \eta^{16} + 6 S_7 \eta^{17} - 20 S_7 \eta^{18} \\
&+8 S_7 \eta^{19} + 20 S_7 \eta^{20} + S_7 \eta^{21} + S_7 \eta^{22} + S_7 \eta^{23} + S_7 \eta^{24} + 134/3 S_7 \eta^{25} \\
&-12 S_7 \eta^{26} + 12 S_7 \eta^{27} + 6 S_7 \eta^{28} - 22/3 S_7 \eta^{29} + 1447/18 S_7 \eta^{30} - 16 S_7 \eta^{31} \\
&+104/3 S_7 \eta^{32} + 6 S_7 \eta^{33} + 201 S_7 \eta^{34} + 2 S_7 \eta^{35} + 25 S_7 \eta^{36} + 56/3 S_7 \eta^{37} + 12 S_7 \eta^{38} \\
&+50 S_7 \eta^{39} + 46 S_7 \eta^{40} + 18 S_7 \eta^{41} + 2 S_7 \eta^{42} + 6 S_7 \eta^{43} + 134/3 S_7 \eta^{44} - 12 S_7 \eta^{45} \\
&+6 S_7 \eta^{46} + 22/3 S_7 \eta^{47} + 6 S_7 \eta^{48} + 134/3 S_7 \eta^{49} - 12 S_7 \eta^{50} + 6 S_7 \eta^{51}.
\end{align*}

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\[-13S_{2,2,2}+6S_{2,2,3}-44/3S_{2,3,2}+38S_{2,3,3}+36S_{2,4,1}+307/6S_{3,1,1}+20S_{3,1,1}D_1^2
\]
\[+86/3S_{3,1,1}\eta+16S_{3,1,2}\eta^2-43/3S_{3,1,2}+10S_{3,1,2}\eta+14S_{3,1,3}-43/3S_{3,2,1}+10S_{3,2,1}\eta
\]
\[+24S_{3,2,2}+22S_{3,3,1}-37/3S_{3,1,1}+26S_{4,1,2}+28S_{4,2,1}+44S_{5,1,1}+40S_{1,1,4}
\]
\[-16/3S_{1,1,3}+16S_{1,1,3}\eta-32S_{1,1,3,2}-24S_{1,1,4,1}-12S_{1,2,2,2}-28/3S_{1,3,1,1}-16S_{3,1,1}\eta
\]
\[-20S_{3,1,2}-20S_{3,2,3}-52S_{4,1,1,1}-12S_{2,1,2,2}-12S_{2,2,2,1}-36S_{3,2,3,1}
\]
\[-12S_{3,1,1,2}-12S_{3,1,2,1}-12S_{3,2,1,1}-12S_{4,1,1,1}-32S_{3,1,1,3}+32S_{1,1,1,1}\right) \quad (3.6)
\]

and

\[
\gamma_{L,1}^{(3)}(N) = 
16 \left( -353/48+119917/864D_1^2+15689/324D_1^3+433/72D_1^4+19/3D_1^5+5/3D_1^6 
-112979/2592\eta+13405/648\eta^2-8045/1296\eta^3 -61/18\eta^4 -1/2\eta^5 -5/6\eta^6 -5\zeta_5 
-15\zeta_5 + 10\zeta_5^2 - 33/8\zeta_4 - 11/4\zeta_4^2 + 235/16\zeta_3 + 8/3\zeta_3^2 + 2\zeta_3 D_1 + 83/8\zeta_3 \eta
+3/2\zeta_3 \eta^2 + 2/3\zeta_3 \eta^3 - 39883/1296S_1 + 19009/324S_1 D_1^2 + 77/9S_1 D_1^3 + 79/6S_1 D_1^4
+20/3S_1 D_1^5 - 19927/324S_1 \eta + 1453/81S_1 \eta^2 - 7/24S_1 \eta^3 + 38/9S_1 \eta^4 + 3S_1 \eta^5
+10S_1 \zeta_5 + 11/2S_1 \zeta_4 - 317/12S_1 \zeta_3 + 4S_1 \zeta_3 D_1^2 + 8/3S_1 \zeta_3 \eta - 8/3S_1 \zeta_3 \eta^2 - 45S_1 \zeta_3 \eta^3
+85175/2592S_2 - 1873/54S_2 D_1^2 - 20/3S_2 D_1^3 - 4S_2 D_1^4 + 4943/648S_2 \eta + 95/216S_2 \eta^2
+229/36S_2 \eta^3 + 25/6S_2 \eta^4 + 2/3S_2 \zeta_2 + 2S_2 \zeta_3 \eta - 4S_2 \zeta_3 \eta^2 - 22247/648S_3 + 241/18S_3 D_1^2
+2/3S_3 D_1^3 - 113/54S_3 \eta + 37/18S_3 \eta^2 + 19/6S_3 \eta^3 - 8/3S_3 \zeta_3 + 725/24S_4 - 16/3S_4 D_1^2
-73/36S_4 \eta + 5/3S_4 \eta^2 - 46/3S_5 + 8/3S_5 \eta + 20/3S_6 + 8/3S_1 D_1^2 + 8/3S_1 D_1^3 + 8S_1 D_1^4
+4/3S_1 \eta + 9/4S_1 \eta^2 - 12S_1 \eta^3 - 14/3S_1 \eta^4 - 8S_1 \zeta_3 \eta + 8S_1 \zeta_3 \eta^2 + 6673/324S_12
-28/3S_1 D_1^2 + 28/3S_1 \eta - 8/3S_1 \eta^2 - 8/3S_1 \eta^3 + 4S_1 \zeta_3 - 605/54S_1,3 + 4/3S_1,3 D_1^2
-14/3S_1,3 \eta - 1/3S_1,3 \eta^2 + 2S_1,3 \eta^3 + 181/18S_1,4 + 10/3S_1,4 \eta - 16/3S_1,5 + 6673/324S_2,1
-28/3S_2,1 D_1^2 + 28/3S_2,1 \eta - 8/3S_2,1 \eta^2 - 8/3S_2,1 \eta^3 + 4S_2,1 \zeta_3 - 1021/54S_2,2 + 4S_2,2 D_1^2
+2/3S_2,2 \eta - 2S_2,2 \eta^2 + 181/18S_2,3 - 8/3S_2,2 \eta + 2S_2,2 \eta^2 - 2S_2,4 - 479/18S_3,1
+16/3S_3,1 D_1^2 + 59/9S_3,1 \eta - 2S_3,1 \eta^2 - 2S_3,1 \eta^3 + 275/18S_3,2 - 10/3S_3,2 \eta - 22/3S_3,3
+43/18S_3,4 - 4/3S_3,4 \eta - 2S_3,4 \eta^2 - 12S_3,4 - 12S_3,5 + 4S_3,1 \eta - 4S_1,1,3 \eta^2 - 20/3S_3,1,4
-20/3S_3,2,3 + 9/4S_3,1,3 - 20/3S_3,1,3 \eta + 4S_3,1,3 \eta^2 + 20/3S_3,1,4 + 20/3S_3,1,4
-20/3S_3,2,3 + 4S_3,2,1 - 20/3S_3,2,1 + 4S_3,2,2 + 8/3S_3,3,1 - 20/3S_3,1,1 + 4/3S_3,1,1 \eta
+16/3S_3,1,2 + 16/3S_3,2,1 + 28/3S_4,1,1 + 16/3S_1,1,3,1 - 8/3S_3,1,1 \right) . \quad (3.7)
\]

The large-\(N\) limit of eq. (3.5) is of the form (2.17) with the large-\(n_c\) cusp anomalous dimension

\[
A_{L,4} = C_F n_c^3 \left( \frac{84278}{81} - \frac{88832}{81} \zeta_2 + \frac{20992}{27} \zeta_3 + 1804\zeta_4 - \frac{352}{3} \zeta_2 \zeta_3 - 352\zeta_5 - 32\zeta_5^2 - 876\zeta_6 \right)
- C_F n_c^2 \eta / \left( \frac{39883}{81} - \frac{26692}{27} \zeta_2 + \frac{16252}{27} \zeta_3 + \frac{440}{3} \zeta_4 - \frac{256}{3} \zeta_2 \zeta_3 - 224 \zeta_5 \right)
\]
Our result for the (complete) $n_f^2$ part was first presented at Loops & Legs 2016, see ref. [31], the rest in a Zurich seminar by one of us [103]. Eq. (3.8) agrees with results of refs. [46, 47], where this quantity was obtained by computing the photon-quark form factor in the large-$n_c$ limit. The lower-order coefficients can be found in eq. (3.11) of ref. [14].

The one- to three-loop coefficients $B_{1,2,3}$ in eq. (4.9) can be found, as coefficients of $\delta(1-x)$, in eqs. (4.5), (4.6) and (4.9) of ref. [14]. The four-loop coefficient in the large-$n_c$ limit reads

\[
B_{L,4} = C_F n_f^2 \left( -\frac{1379569}{5184} + \frac{24211}{27} \zeta_2 - \frac{9803}{162} \zeta_3 - \frac{9382}{9} \zeta_4 + \frac{838}{9} \zeta_2 \zeta_3 + 1002 \zeta_5 \\
+ \frac{16}{3} \zeta_3^2 + 135 \zeta_6 - 80 \zeta_2 \zeta_5 + 32 \zeta_3 \zeta_4 - 560 \zeta_7 \right) \\
+ C_F n_c^2 n_f \left( \frac{353}{3} - \frac{85175}{162} \zeta_2 - \frac{137}{9} \zeta_3 + \frac{16186}{27} \zeta_4 - \frac{584}{9} \zeta_2 \zeta_3 - \frac{248}{3} \zeta_5 - \frac{16}{3} \zeta_3^2 - 144 \zeta_6 \right) \\
- C_F n_c n_f \left( \frac{127}{18} - \frac{5036}{27} \zeta_2 + \frac{932}{27} \zeta_3 + \frac{1292}{27} \zeta_4 - \frac{160}{9} \zeta_2 \zeta_3 - \frac{32}{3} \zeta_5 \right) \\
- C_F n_c^2 \left( \frac{131}{18} - \frac{32}{81} \zeta_2 - \frac{304}{81} \zeta_3 + \frac{32}{27} \zeta_4 \right). \tag{3.9}
\]

The coefficients $B_n$ contain collinear contributions to the evolution kernels. With the help of the QCD corrections to the quark form factor in dimensional regularization, one can extract from them the universal eikonal anomalous dimension. The latter governs the subleading infrared poles in gauge-theory amplitudes and captures contributions from large-angle soft gluons [100, 104, 105].

As mentioned above, the coefficients $C_n$ and $\tilde{D}_n$ in eq. (4.9) do not provide new information, but are functions of lower-order quantities. They are given by

\[
C(a_s) = (A(a_s))^2, \quad \tilde{D}(a_s) = A(a_s) \cdot (B(a_s) - \beta(a_s)/a_s), \tag{3.10}
\]

cf. ref. [65], which leads to the four-loop relations

\[
C_4 = A_2^2 + 2 A_1 A_3, \quad \tilde{D}_4 = \sum_{k=1}^{3} A_k \cdot (B_{4-k} - \beta_{3-k}). \tag{3.11}
\]

Using the results (3.8) and (3.9), it is actually now possible to predict $C_5$ and $\tilde{D}_5$ for large $n_c$.

The new functions (3.6) and (3.7) are shown in figures 2 and 3, respectively, together with their large-$N$ approximation (2.17) with the coefficients given above. In the right panels, the results are divided by $\ln N$, so for $N \to \infty$ the curves tend to constants given by the respective terms in the four-loop cusp anomalous dimension (3.8).

The approach to this asymptotic behaviour is very slow: the $n_f^0$ contribution in figure 2 is 0.856 of its asymptotic result at $N = 30$, yet it deviates by less than 10% only from an $N$-value above $N = 500$. The corresponding numbers for the $n_f^1$ part in figure 3 are
Figure 2. The $n_f^0$ part (3.6) of the anomalous dimensions $\gamma_{ns}^{(n)\pm}(N)$ in the large-$n_c$ limit, compared with its large-$N$ expansion with all terms included in eq. (2.17) and, in the right panel, its asymptotic behaviour for $N \to \infty$. The exact curve has been computed by via the $x$-space counterpart of eq. (3.6), see section 4.

0.873 at $N = 30$ and $N \simeq 185$ for a deviation by less than 10%. It might be interesting to note, on the other hand, that the corresponding $n_f^a$ coefficient of $A_n$, here and in all lower-order cases (in full QCD), falls in the interval spanned by the corresponding results for $\gamma_{ns}^{(n)}(2)/\ln 2$ and $\gamma_{ns}^{(n)}(4)/\ln 4$.

The results (3.1) for $\gamma_{ns}^{(3)+}$ (closed circles) and (3.2) for $\gamma_{ns}^{(3)-}$ (open circles) are shown for the physically relevant values of $n_f$ in figure 4, together with the all-$N$ results in the large-$n_c$ limit. As at the previous orders in $\alpha_s$, there are cancellations between the $n_f$-independent and the $n_f$-dependent contributions, which are particularly pronounced here at $n_f = 5$. For this number of light flavours, which is relevant for high-energy processes at the LHC, the large-$n_c$ result do not describe the (small) fourth-order QCD contributions to the non-singlet evolution equations at the phenomenologically most relevant moments $N$ and momentum fractions $x$. We therefore need to convert the calculated moments to practically usable constraints on the four-loop splitting functions $P_{ns}^{(3)}(x)$.

4 Results in $x$-space

The four-loop non-singlet splitting functions $P_{ns}^{(3)}(x)$ are derived from the all-$N$ results for the corresponding anomalous dimensions by an inverse Mellin transformation that expresses these functions in terms of harmonic polylogarithms (HPLs). This transformation can be
Figure 3. As figure 2, but for the $n_f^3$ contribution (3.7). The value of $\gamma_{ns}^{(n_f)}(N)/\ln N$ in the limit $N \to \infty$ is given by the corresponding coefficient of $A_{L,4}$ in the second line of eq. (3.8).

Figure 4. Our even-$N$ results for $\gamma_{ns}^{(n_f)+}(N)$ and odd-$N$ values for $\gamma_{ns}^{(n_f)-}(N)$ at $n_f = 3, \ldots, 6$, compared with their common large-$n_c$ limit now known at all $N$. The results have been converted to an expansion in $\alpha_s$. 
performed by an algebraic procedure \[87, \ 88\] based on the fact that harmonic sums occur as coefficients of the Taylor expansion of HPLs.

For the convenience of the reader, we recall their basic definitions \[87\]. The lowest-weight \((w = 1)\) functions \(H_w(x)\) are given by

\[
H_0(x) = \ln x \, , \quad H_{\pm 1}(x) = \mp \ln(1 \mp x) \, . \tag{4.1}
\]

The higher-weight \((w \geq 2)\) functions are recursively defined as

\[
H_{m_1, \ldots, m_w}(x) = \begin{cases} 
\frac{1}{w!} \ln^w x \, , & \text{if } m_1, \ldots, m_w = 0, \ldots, 0 \\
\int_0^x dz \, f_{m_1}(z) H_{m_2, \ldots, m_w}(z) \, , & \text{otherwise} 
\end{cases} \tag{4.2}
\]

with

\[
f_0(x) = x^{-1} \, , \quad f_{\pm 1}(x) = (1 \mp x)^{-1} \, . \tag{4.3}
\]

For chains of indices 'zero' we employ the abbreviated notation

\[
H_{0, \pm 1, 0, \ldots, 0, \pm 1, \ldots}(x) = H_{\pm(m+1), \pm(n+1), \ldots}(x) \, . \tag{4.4}
\]

The argument \(x\) will be suppressed in all results below, and we express the terms with \((1 \pm x)^{-1}\) in terms of the \(x\)-dependence of the leading-order splitting function \(P_{qq}^{(0)}\),

\[
p_{qq}(x) = 2(1 - x)^{-1} - 1 - x \, . \tag{4.5}
\]

In this notation, the common large-\(n_c\) limit of the functions \(P_{ns}^{(3)+}(x)\) and \(P_{ns}^{(3)-}(x)\) is given by

\[
P_{ns, L}(x) = C_F \left( n_c^2 P_{L, 0}^{(3)}(x) + n_c n_f P_{L, 1}^{(3)}(x) + n_c n_f^2 P_{L, 2}^{(3)}(x) + n_f^3 P_{L, 3}^{(3)}(x) \right) \tag{4.6}
\]

with the new results

\[
P_{L, 0}^{(3)}(x) = \\
p_{qq}(x) \left( \frac{42139}{81} - 160 \zeta_6 - 204 \zeta_5 + \frac{7666}{9} \zeta_4 + \frac{10334}{27} \zeta_3 + 16 \zeta_2 - \frac{44416}{81} \zeta_2 \zeta_3 \right.
\]

\[
+ \frac{71591}{108} H_{0, 0, 0} + 504 H_{0, 0, 3} + \frac{212}{3} H_{0, 0, 0} \zeta_4 - \frac{6899}{9} H_{0, 0, 0} \zeta_2 + \frac{44416}{81} H_{1, 0} + 680 H_{1, 0} \zeta_4
\]

\[
+ 128 H_{1, 0} \zeta_3 - \frac{5788}{9} H_{1, 0} \zeta_2 + 512 H_{1, 1} \zeta_4 + \frac{128}{3} H_{1, 1} \zeta_3 + 192 H_{1, 1} \zeta_3 + \frac{5788}{9} H_{1, 3}
\]

\[
- 160 H_{1, 3} \zeta_2 + \frac{448}{3} H_{2, 1, 0} + 368 H_{1, 4} + \frac{9029}{27} H_{2, 0} + 32 H_{2, 0} \zeta_3 - \frac{352}{3} H_{2, 0} \zeta_2 + 192 H_{2, 1} \zeta_3
\]

\[
+ \frac{1072}{3} H_{2, 2} - 96 H_{2, 2} \zeta_2 + \frac{352}{3} H_{2, 2} + 288 H_{2, 4} + \frac{5581}{9} H_{3, 0} - 176 H_{3, 0} \zeta_2 + \frac{1228}{3} H_{3, 1}
\]

\[
- 96 H_{3, 1} \zeta_2 + \frac{344}{3} H_{3, 2} + 176 H_{3, 3} + 228 H_{4, 0} + \frac{296}{3} H_{4, 1} + 224 H_{4, 2} + 288 H_{5, 0} + 352 H_{5, 1}
\]

\[
+ \frac{55291}{108} H_{0, 0, 3} - 32 H_{0, 0, 0} \zeta_3 - 252 H_{0, 0, 0} \zeta_2 + \frac{7120}{27} H_{1, 0} - 16 H_{1, 0} \zeta_3 - \frac{448}{3} H_{1, 0} \zeta_2
\]
\[+64H_{1,1,0}\zeta_3 + \frac{128}{3} H_{1,1,0}\zeta_2 + 256H_{1,1,1}\zeta_3 - \frac{128}{3} H_{1,1,1} + 192H_{1,1,4} + \frac{1072}{3} H_{1,2,0} + 304 \frac{H_{1,3,0} - \frac{224}{3} H_{1,3,1} + 160H_{1,3,2} + 400H_{1,4,0} + 416H_{1,4,1} + \frac{3635}{9} H_{2,0,0}}{3} - 288H_{2,0,0}\zeta_2 + \frac{1072}{3} H_{2,1,0} + 104H_{2,2,0} + 96H_{2,2,2} + 304H_{2,3,0} + 288H_{2,3,1} + 572 \frac{H_{3,0,0} + \frac{344}{3} H_{3,1,0} + 96H_{3,1,2} + 192H_{3,2,0} + 96H_{3,2,1} + 208H_{4,0,0} + 224H_{4,1,0}}{3} + 96H_{4,1,1} + 545H_{0,0,0,0,0} - 224H_{0,0,0,0,0}\zeta_2 + \frac{4745}{9} H_{1,0,0,0,0} - 368H_{1,0,0,0,0}\zeta_2 - 192H_{1,0,0,0,0}\zeta_2 + 256H_{1,1,0}\zeta_2 - 256H_{1,1,1} + 256H_{1,1,3} + 256H_{1,1,3,0} + 256H_{1,1,3,1} + \frac{176}{3} H_{1,2,0,0} + 96H_{1,2,2,0} + 176H_{1,3,0,0} + 160H_{1,3,1,0} + 124H_{2,0,0,0} + \frac{176}{3} H_{2,1,0,0} + 96H_{2,1,2,0} + 96H_{2,2,1,0} + 112H_{3,0,0,0} + 112H_{3,1,0,0} + 96H_{3,1,1,0} + \frac{520}{3} H_{0,0,0,0,0} + \frac{560}{3} H_{1,0,0,0,0} - \frac{160}{3} H_{1,1,0,0,0} - 48H_{1,2,0,0,0} + 128H_{2,1,0,0,0} + 80H_{3,0,0,0,0} + 160H_{1,0,0,0,0,0} + 64H_{1,1,0,0,0,0} - 320H_{1,1,1,0,0,0} + \frac{24211}{54} H_{0} - 160H_{0} + 538 \frac{H_{0}\zeta_4}{3} + \frac{193}{3} H_{0}\zeta_3 - \frac{3700}{9} H_{0}\zeta_2 + 96H_{0}\zeta_2 - 80H_{1}\zeta_3 - \frac{8}{3} H_{1}\zeta_4 + \frac{644}{9} H_{1}\zeta_3 + 96H_{1}\zeta_2 + \frac{44416}{81} H_{2} + 552H_{2}\zeta_3 + \frac{272}{3} H_{2}\zeta_3 - \frac{1072}{3} H_{2}\zeta_2 + \frac{3700}{9} H_{3} + 80H_{3}\zeta_3 - \frac{344}{3} H_{3}\zeta_2 \times H_{4} - 224H_{4}\zeta_2 + 252H_{5} + 224H_{6} + x \left( - \frac{184}{3} \zeta_5 - \frac{80}{3} \zeta_2 \zeta_3 + \frac{796}{3} \zeta_4 - \frac{871}{3} \zeta_3 \right) + \frac{6835}{54} \zeta_2 - \frac{567245}{324} H_{0,0,0,0,0} - \frac{52}{3} H_{0,0,0,0,0} + \frac{875}{3} H_{0,0,0,0,0} + \frac{1015}{3} H_{0,0,0,0,0} - \frac{944}{3} H_{0,0,0,0,0} + \frac{254}{3} H_{1,1,1} - \frac{116}{3} H_{1,2,2} + \frac{944}{3} H_{1,2,2} - \frac{1130}{3} H_{1,3} + \frac{40}{3} H_{1,3} - 96H_{3,2} - \frac{1060}{3} H_{1,4,0} - 216H_{1,1,1} - \frac{33091}{36} H_{0,0,0,0,0} + 428H_{0,0,0,0,0} + \frac{1223}{3} H_{0,0,0,0,0} - \frac{116}{3} H_{0,0,0,0,0} - 56H_{2,2,2,0} - 148H_{3,0,0,0,0} - 96H_{3,1,0,0,0,0} - 550H_{0,0,0,0,0} + \frac{152}{3} H_{0,0,0,0,0} - \frac{952}{3} H_{0,0,0,0,0} - 280H_{0,0,0,0,0} - \frac{237016}{81} H_{0,0,0,0,0} - \frac{422}{3} H_{0,0,0,0,0} + \frac{569}{3} H_{0,0,0,0,0} + \frac{937}{3} H_{0,0,0,0,0} - \frac{6435}{54} H_{0,0,0,0,0} + \frac{116}{3} H_{0,0,0,0,0} - \frac{937}{3} H_{0,0,0,0,0} + \frac{96}{3} H_{3,2} - \frac{875}{3} H_{4,4,0,0,0,0} + \left( 1 + x \right) \left( 360H_{0,0,0,0,0} + 64H_{0,0,0,0,0} + 268H_{0,0,0,0,0} + 80H_{0,0,0,0,0} \right) - 160H_{2,1,1,0,0,0,0} + 112H_{3,0,0,0,0,0} - 112H_{3,0,0,0,0,0} - 88H_{4,2,0,0,0} - 172H_{5,0,0,0,0} - 176H_{5,0,0,0,0} + 24H_{0,0,0,0,0,0} + 112H_{2,0,0,0,0} + 96H_{2,0,0,0,0} + 48H_{2,0,0,0,0} - 32H_{2,0,0,0,0} - 48H_{3,0,0,0,0} - 100H_{4,1,0,0,0} + 168H_{0,0,0,0,0} + 64H_{2,2,0,0,0} - 24H_{3,0,0,0,0,0} + 64H_{3,1,0,0,0,0} - 16H_{2,1,0,0,0,0} + 128H_{2,1,0,0,0,0} - 70H_{0,0,0,0,0,0} + 40H_{0,0,0,0} - 48H_{0,0,0,0} + 216H_{2,0,0,0,0,0} + 88H_{2,0,0,0,0,0} - 168H_{6} \left( 1 + x \right) \left( \frac{1325880}{324} + \frac{43301}{27} \right) H_{1,0} + \frac{256}{3} H_{1,0,0,0,0,0} + \frac{3224}{3} H_{1,1,0,0,0,0} + \frac{1448}{3} H_{1,1,0,0,0,0} + \frac{560}{3} H_{1,1,0,0,0,0} - 560 H_{1,1,0,0,0,0} - 192H_{1,1,0,0,0,0}
\[ P_{L,1}^{(3)}(x) = \]
\[ p_{eq}(x) \left( -\frac{39883}{162} + \frac{128}{3} \delta_0 - \frac{680}{9} \zeta_4 - \frac{8126}{27} \zeta_3 + \frac{13346}{81} \zeta_2 + \frac{32}{3} \zeta_2 \zeta_3 - \frac{22247}{81} \delta_0 \right) + \]
\[ \frac{128}{3} H_{0,0,0,0} \zeta_3 + \frac{1372}{9} H_{0,0,0,1} - \frac{13346}{81} H_{1,0,0,3} + \frac{752}{9} H_{1,0,2,0} - \frac{128}{3} H_{1,1,3} + \]
\[ \frac{752}{9} H_{1,1,4} - \frac{4084}{27} H_{2,0} + \frac{64}{3} H_{2,0} \zeta_2 - \frac{160}{3} H_{2,2} - \frac{64}{3} H_{2,3} - \frac{1100}{9} H_{3,0} + \]
\[ \frac{160}{3} H_{3,1,1} - \frac{128}{3} H_{3,2,0} - \frac{964,0}{3} H_{4,1} - \frac{725}{3} H_{0,0,0,2} + \frac{128}{3} H_{1,0,0,2} + \frac{32}{3} H_{1,0,0,3} - \frac{128}{3} H_{1,1,0,3} - \frac{160}{3} H_{1,2,0,0} - \frac{160}{3} H_{1,3,0,0} + \frac{64}{3} H_{3,1,1} - \frac{724}{9} H_{2,0,0,0} + \]
\[ \frac{32}{3} H_{2,1,0,3} - \frac{32}{3} H_{2,2,0,0} - \frac{176}{3} H_{3,0,0,0} - \frac{128}{3} H_{3,1,0,0} + \frac{368}{3} H_{0,0,0,0} - \frac{724}{9} H_{1,0,0,0} - \frac{32}{3} H_{2,1,0,0} - \frac{160}{3} H_{2,2,0,0,0} - \frac{128}{3} H_{2,3,0,0,0} + \frac{32}{3} H_{2,4,0,0,0} - \frac{128}{3} H_{1,0,0,0,0} + \frac{160}{3} H_{1,1,0,0,0} - \frac{85175}{324} H_{0,0} - \frac{756}{3} H_{0,0,0,3} + \frac{1916}{9} H_{0,0,0,2} - \frac{208}{3} H_{1,0,0,2} - \frac{208}{3} H_{0,0,0,0,0} - \frac{13346}{81} H_{2} - \frac{224}{3} H_{2,0} - \frac{160}{3} H_{2,2} - \frac{1916}{9} H_{3} + \frac{128}{3} H_{3,0} - \frac{1372}{9} H_{4} - \frac{964,0}{3} H_{5} \right) + \]
\[ x \left( -\frac{496}{3} \zeta_4 + \frac{454}{3} \zeta_3 - \frac{9752}{27} \zeta_2 + \frac{4128}{81} H_{0,0,0} + \frac{28}{3} H_{0,0,0,2} + \frac{170}{3} H_{2,0,0} + \frac{320}{3} H_{2,0,0} \right) - \frac{92}{3} H_{2,1,0} - \frac{32}{3} H_{2,2} - \frac{32}{3} H_{2,3} + \frac{224}{3} H_{3,0} + \frac{32}{3} H_{3,1} + \frac{2339}{9} H_{0,0,0,0} + \frac{428}{3} H_{2,0,0} + \frac{32}{3} H_{2,1,0} + \frac{376}{3} H_{0,0,0,0} + \frac{304}{3} H_{2,1,0,0} + \frac{115273}{108} H_{0} - \frac{236}{3} H_{0,0,0} - \frac{74}{3} H_{0,0} - \frac{9752}{27} H_{2}, \]
\[-\frac{32}{3} H_2 \zeta_2 + \frac{74}{3} H_3 - \frac{28}{3} H_4 \] 
\[-\left(1 + x\right) \left(-124\zeta_6 + \frac{88}{3} \zeta_5 - 32\zeta_3^2 - \frac{16}{3} \zeta_2 \zeta_4 + \frac{64}{3} H_0 \zeta_3 \right) \]
\[+ 32 H_2 \zeta_3 + 64 H_2,1 \zeta_3 + 32 H_{2,4} - 32 H_{3,0} \zeta_2 + 32 H_{3,3} + \frac{88}{3} H_4,0 - 48 H_{0,0,0} \zeta_2 \]
\[+ 20 H_{0,0,0,0,0} + 16 H_0 \zeta_4 + 64 H_2 \zeta_4 + 32 H_3 \zeta_3 + 48 H_5 \] 
\[+ (1 - x) \left(-\frac{136319}{81} - \frac{12148}{27} H_{1,0} + 96 H_{1,0} \zeta_3 - \frac{320}{3} H_{1,0} \zeta_2 - \frac{872}{3} H_{1,1,1} + 192 H_{1,1} \zeta_3 - \frac{64}{3} H_{1,2} + \frac{320}{3} H_{1,1} \right) \]
\[+ 96 H_{1,4} - \frac{1036}{9} H_{1,0,0} - 96 H_{1,0,0} \zeta_2 - \frac{64}{3} H_{1,1,0} - 192 H_{1,1,0} \zeta_2 + 192 H_{1,1,3} - \frac{64}{3} H_{1,2,0} \]
\[+ \frac{16}{3} H_{1,0,0,0} - 128 H_{1,1,0,0} - 96 H_{1,2,0,0} - 192 H_{1,1,1,0,0} - \frac{11416}{27} H_{1} + 192 H_1 \zeta_4 \]
\[+ \frac{320}{3} H_1 \zeta_2 + \frac{64}{3} H_1 \zeta_3 + 96 H_2 \zeta_3 \]
\[+ \frac{224}{3} \zeta_4 - 38 \zeta_3 + \frac{1616}{3} \zeta_2 - \frac{38036}{81} H_0,0 + 20 H_{0,0} \zeta_2 \]
\[+ \frac{1954}{9} H_2,0 - \frac{256}{3} H_{2,0} \zeta_2 - \frac{404}{3} H_{2,1} - 32 H_{2,2} + \frac{256}{3} H_{2,3} - \frac{128}{3} H_{3,0} - \frac{160}{3} H_{3,1} \]
\[+ \frac{473}{9} H_{0,0,0} - \frac{116}{3} H_{0,0,0} - 32 H_{2,1,0} + \frac{152}{3} H_{0,0,0,0} - \frac{272}{3} H_{2,1,0,0} - \frac{34035}{324} \frac{3}{3} H_0 \]
\[+ \frac{44}{3} H_0 \zeta_3 + \frac{2090}{9} H_0 \zeta_2 - \frac{1616}{3} H_2 + 32 H_2 \zeta_2 - \frac{2090}{9} H_3 - 20 H_4 + \delta(1 - x) \left(\frac{353}{3} \right) \]
\[+ \frac{85175}{162} \zeta_2 - \frac{137}{9} \zeta_3 + \frac{16186}{27} \zeta_4 - \frac{584}{9} \zeta_2 \zeta_3 - \frac{248}{3} \zeta_5 - \frac{16}{3} \zeta_3^2 - 144 \zeta_6 \right). \tag{4.8} \]

The $n_2^3$ and $n_3^3$ contributions to eq. (4.6) are given by eqs. (4.6) and (4.12) of ref. [32].

Disregarding terms that vanish for $x \to 1$, the large-$x$ behaviour of $P_{n_s L}^{(3)}(x)$ can be written as

\[ P_{n_s L}^{(n-1)}(x) = \frac{A_{L,n}}{(1 - x)_+} + B_{L,n} \delta(1 - x) + C_{L,n} \ln(1 - x) - A_{L,n} + \tilde{D}_{L,n} \] \tag{4.9} 

in terms of the coefficients specified in eqs. (3.8), (3.9) and (3.11) above. The numerical values of the coefficients of the small-$x$ logarithms, $\ln^k x$ with $k = 1, \ldots, 6$, can be read off from eq. (4.11) below. All six logarithms and the constant contribution for $x \to 0$ are required for a good approximation to the splitting functions at small $x$-values relevant for collider physics.

In view of the length and complexity of the exact expressions (4.7) and (4.8) it is useful to have at one’s disposal also compact approximate representations involving, besides powers of $x$, only simple functions like the plus-distribution and the end-point logarithms

\[ D_0 = 1/(1 - x)_+, \quad x_1 = 1 - x, \quad L_1 = \ln(1 - x), \quad L_0 = \ln x. \] \tag{4.10} 

Such approximations can be readily used in $N$-space evolution programmes, see, e.g., ref. [106]. The results (4.7) and (4.8) can be parametrized with a high accuracy (of 0.1%
or better) as

$$\begin{align*}
C_F n_c^3 P_{L,0}^{(3)}(x) + C_F n_c^2 n_f P_{L,1}^{(3)}(x) &= \\
&= 21209.02 D_0 + (25796.09 - 1.0) \delta(x_1) + 19069.80 L_1 - 29733.85 \\
&+ 25000 \left( x_1 (3.5242 + 8.3679 x - 1.2395 x^2 + 1.7423 x^3) + 11.916 x L_0 \\
&\quad - 0.2237 x L_0^2 - 0.0129 x L_0^3 + 11.937 x_1 L_1 + 13.955 L_0 L_1 \right) \\
&+ 516713.33 L_0 + 17120.95 L_0^2 + 2863.226 L_0^3 + 297.8255 L_0^4 + 16 L_5^5 + 1/2 L_6^6 \\
&+ n_f \left( -5179.372 D_0 - (5818.637 + 0.35) \delta(x_1) - 3079.761 L_1 + 8115.605 \\
&\quad + 2500 \left( x_1 (-7.4077 + 4.5141 x - 1.0069 x^2 + 0.7641 x^3) + 8.4211 x_1 L_1 \\
&\quad + 7.5633 L_0 L_1 + 7.5236 x L_0 + 0.2208 x L_0^2 + 0.05712 x L_0^3 \right) \\
&\quad - 9239.374 L_0 - 2917.312 L_0^2 - 430.5308 L_0^3 - 36 L_0^4 - 4/3 L_0^5 \right) .
\end{align*}$$

(4.11)

Here the exact large-\(x\) and small-\(x\) coefficients have been rounded to seven significant figures. The brackets multiplied by 25000 and 2500 have been obtained by fits to the exact expressions at \(10^{-6} \leq x \leq 1 - 10^{-6}\). The small shifts of the \(\delta(1-x)\) fine-tune the accuracy of the resulting low moments and of the convolutions with the quark distributions. The required evaluation of the HPLs has been performed using a weight-6 extension of ref. [107] and the program of ref. [108], which return identical results at the accuracy considered here.

For the corresponding non-leading contributions in the large-\(n_c\) limit, denoted by the subscript \(N\) in eqs. (3.1) and (3.2) above, we are for now limited to approximations analogous to (but more accurate than) those once constructed at three loops [48–50]. For the \(n_f^0\) and \(n_f^1\) parts of \(P_{ns}^{(3)+}(x)\) we employ an ansatz consisting of

- the two large-\(x\) parameters \(A_4\) and \(B_4\), cf. eqs. (2.17) and (4.9),
- two of three suppressed large-\(x\) logarithms \((1-x) \ln^k(1-x), k = 1, 2, 3,\)
- one of ten two-parameter polynomials in \(x\) that vanish for \(x \to 1,\)
- two of the three unknown small-\(x\) logarithms \(\ln^k x, k = 1, 2, 3.\)

The parameters of the 90 resulting trial functions are determined from the eight available moments, and then two representatives as chosen that indicate the remaining uncertainty. The result of this process is illustrated in figures 5 and 6.

Supplementing the approximations \(A\) and \(B\) in figures 5 and 6 by accurate parametrizations of the complete \(n_f^2\) results of ref. [32] and the exact (but numerically truncated) \(n_f^3\) expressions in a non-HPL notation, we obtain

$$\begin{align*}
P_{ns,A}^{(3)+}(x) &= C_F n_c^3 P_{L,0}^{(3)}(x) + C_F n_c^2 n_f P_{L,1}^{(3)}(x) + P_{L,n_f}^{(3)+}(x) \\
&= -507.152 D_0 - 2405.03 \delta(x_1) - 1777.27 L_1^2 x_1 - 204.183 L_1^3 x_1 \\
&\quad + 3948.16 x_1 - 2464.61 (2x - x^2) x_1 - 1839.44 L_0^4 - 402.156 L_0^3
\end{align*}$$
Figure 5. About 90 trial functions for the $n_f$-independent contribution to the non-leading ($N$) large-$n_c$ part of splitting function $P_{n_s}^{(3)+}(x)$, multiplied by $x^{0.4}(1-x)$ for display purposes. The two functions chosen to represent the remaining uncertainty are denoted by $A$ and $B$ and shown by solid (blue) lines.

\[
- 55.87553 L_0^4 - 2.831276 L_0^5 - 0.1488340 L_0^6 - 2601.749 - 2118.867 L_1 \\
+ n_f \left( 7.33927 (D_0 - 1) + 267.965 \delta(x_1) - 143.813 L_1 x_1 - 18.8803 L_1^2 x_1 \\
- (1116.34 - 1071.24 x) x_1 - 59.3041 L_0^3 - 8.4620 L_0^3 \\
+ 4.658436 L_0^4 + 0.2798354 L_0^5 + 312.1643 + 337.9310 L_1 \right) \tag{4.12}
\]

and

\[
P_{n_s,B}^{(3)+}(x) = C_{F} n_c^3 P_{L,0}^{(3)}(x) + C_{F} n_c^2 n_f P_{L,1}^{(3)}(x) + P_{L_{n_f}}^{(3)+}(x) \\
- 505.209 (D_0 - 1) - 2394.47 \delta(x_1) - 173.936 L_1^2 x_1 + 223.078 L_1^3 x_1 \\
+ (8698.39 - 10490.47 x) x_1 + 1389.73 L_0 + 189.576 L_0^2 \\
- 55.87553 L_0^4 - 2.831276 L_0^5 - 0.1488340 L_0^6 - 2601.749 - 2118.867 L_1 \\
+ n_f \left( 7.53662 (D_0 - 1) + 269.028 \delta(x_1) - 745.573 L_1 x_1 + 8.61438 L_1^2 x_1 \\
- (690.151 + 656.386 x^2) x_1 + 133.702 L_0^2 + 34.0569 L_0^3 \\
+ 4.658437 L_0^4 + 0.2798354 L_0^5 + 312.1643 + 337.9310 L_1 \right) \tag{4.13}
\]
Figure 6. As figure 5, but for the $n_f^1$ contribution. The ratio in the right panel is shown for a smaller $x$-range than in figure 5 due to a sign change of the function at $x \approx 0.3$. The large relative width of the uncertainty band close to $x = 1$ is due to another change of sign at $(1-x) \approx 0.005$.

with

$$P_{L_{n_f}}^{(3)+}(x) =$$

$$n_f^2 \left( 195.5772 D_0 + 26.68861 L_1 - 376.0092 + (193.8554 + 0.0037) \delta(x_1) ight.$$  

$$+ 250 (x_1 (3.0008 + 0.8619 x - 0.12411 x^2 + 0.31595 x^3) - 0.37529 x L_0$$  

$$- 0.21684 x L_0^3 - 0.02295 x L_0^3 + (0.03394 x_1 + 0.40431 L_0) L_1)$$  

$$+ 393.0056 L_0 + 112.5705 L_0^3 + 16.52675 L_0^3 + 0.7901235 L_0^4$$

$$+ n_f^3 \left( 3.272344 D_0 + 3.014982 \delta(x_1) - 2.426296 - 0.8460488 x$$  

$$+ (0.5267490 x_1^{-1} - 3.687243 + 3.160494 x) L_0 - (0.1316872 (10 x_1^{-1} + 1)$$  

$$- 1.448560 x) L_0^2 - (0.2633745 x_1^{-1} - 0.1316870 (1 + x)) L_0^3 \right).$$ (4.14)

The case of $P_{ns}^{(3)-}(x)$ can be treated in the same manner, but taking into account that only its leading small-$x$ logarithm is known up to now [92]. After careful consideration, the two approximations indicating the uncertainty band in this case are chosen as

$$P_{ns,A}^{(3)-}(x) = C_F n_c^3 P_{L,0}^{(3)}(x) + C_F n_c^2 n_f P_{L,1}^{(3)}(x) + P_{L_{n_f}}^{(3)-}(x)$$
\[-511.228 (D_0 - 1) - 2426.05 \delta(x_1) + 31897.82 L_1 x_1 + 4653.76 L_1^2 x_1 \\
+ (5992.88 (1 + 2x) + 31321.44 x^2) x_1 - 1618.07 L_0 + 2.25480 L_0^3 \\
+ 0.4964335(L_0^6 + 6 L_0^5) - 2601.749 - 2118.867 L_1 \\
+ n_f \left( 7.08645 (D_0 - 1) + 266.669 \delta(x_1) + 1856.63 L_1 x_1 + 440.17 L_1^2 x_1 \\
\quad + (114.457 (1 + 2x) + 2570.73 x^2) x_1 - 127.012 L_0^2 + 2.69618 L_0^4 \\
\quad + 312.1643 + 337.9310 L_1 \right) \right) \\
\] (4.15)

and
\[
P_{\text{ns,B}}^{(3)-}(x) = C_F n_c^2 P_{L_0}^{(3)}(x) + C_F n_c^2 n_f P_{L_1}^{(3)}(x) + P_{L_{n_f}}^{(3)-}(x) \\
- 502.481 (D_0 - 1) - 2380.255 \delta(x_1) - 3997.39 L_1 x_1 + 511.567 L_1^2 x_1 \\
+ (4043.59 - 15386.6 x) x_1 + 1532.96 L_0^3 + 31.023 L_0^3 \\
+ 0.4964335 (L_0^6 + 18 L_0^5) - 2601.749 - 2118.867 L_1 \\
+ n_f \left( 7.82077(D_0 - 1) + 270.468 \delta(x_1) - 1360.04 L_1 x_1 + 38.7337 L_1^2 x_1 \\
\quad - (335.995 (2 + x) + 1605.91 x^2) x_1 - 9.76627 L_0^2 + 0.14218 L_0^5 \\
\quad + 312.1643 + 337.9310 L_1 \right) \right) \\
\] (4.16)

with
\[
P_{L_{n_f}}^{(3)-}(x) = \\
n_f^2 \left( 195.5772 D_0 + 26.68861 L_1 - 376.0092 + (193.8554 + 0.0037) \delta(x_1) \\
+ 250 (x_1 (3.2206 + 1.7507 x + 0.13281 x^2 + 0.45969 x^3) + 1.5641 x L_0 \\
- 0.37902 x L_0^2 - 0.03248 x L_0^3 + (2.7511 x_1 + 3.2709 L_0) L_1 \\
+ 437.8810 L_0 + 128.2948 L_0^2 + 19.59945 L_0^3 + 0.9876543 L_0^4 \right) \\
+ n_f^3 \left( 3.272344 D_0 + 3.014982 \delta(x_1) - 2.426296 - 0.8460488 x \\
+ (0.5267490 x_1^{-1} - 3.687243 + 3.160494 x) L_0 - (0.1316872 (10 x_1^{-1} + 1) \\
- 1.448560 x) L_0^2 - (0.2633745 x_1^{-1} - 0.1316870 (1 + x)) L_0^3 \right). \\
\] (4.17)

The \(n_f^3\) contribution to this last equation is the same as in eq. (4.14).

Before we illustrate these results, it is useful to briefly recall the behaviour of the corresponding third-order splitting functions. This is done in figure 7 for \(n_f = 4\) flavours. The corresponding size and uncertainty bands of \(P_{\text{ns}}^{(3)+}(x)\) and \(P_{\text{ns}}^{(3)-}(x)\) are shown in figure 8 together with their large-\(n_c\) limit. The qualitative pattern and the rough size of
Figure 7. The three-loop splitting functions $P_{ns}^{(2)\pm}(x)$ and their large-$n_c$ limit for QCD with four flavours.

Figure 8. The four-flavour uncertainty bands for the four-loop splitting functions $P_{ns}^{(3)\pm}(x)$ generated by eqs. (4.12) and (4.13) for $a = +$ and by eqs. (4.15) and (4.16) for $a = -$, compared to their exact large-$n_c$ limit. As in figure 7, the curves are scaled to an expansion in $\alpha_s$, and only the $x < 1$ contributions are shown.
the corrections as coefficients of $\alpha_s^3$ and $\alpha_s^4$, respectively, are comparable in the region of $x$ for which definite conclusions can be drawn.

The four-loop ‘sea’ contribution $P^{(3)s}_{\text{ns}}(x)$ to the evolution of the total valence distribution is suppressed by two powers of $(1-x)$ for $x \to 1$, but its $n_f^1$ part is completely unknown in the small-$x$ limit. In this case, we use the nine odd moments (3.3) with a suitably modified ansatz, in which the coefficient of $\ln x$ is varied ‘by hand’ over a sufficiently wide range, and the coefficients of $\ln^k x$, $k = 1, \ldots, 5$, are all determined from the moments. In this manner we obtain

$$P^{(3)s}_{\text{ns},A/B}(x) = n_f P^{(3)s}_{1,A/B}(x) + n_f^2 P^{(3)s}_{2}(x)$$  \hspace{1cm} (4.18)

with

$$P^{(3)s}_{1,A}(x) = 60.40 x_1 L_1^4 + 4.685 x_1 L_1^3 + x_1 (4989.2 - 1607.73 x)$$  \hspace{1cm} (4.19)

$$+ 3687.6 L_0^2 + 3296.6 L_0 + 1271.11 L_0^3 + 533.44 L_0^4 + 97.27 L_0^5 + 4 L_0^6,$$

$$P^{(3)s}_{1,B}(x) = -254.63 x_1 L_1 - 0.28953 x_1 L_1^3 + 1030.79 x_1 x + 1266.77 x_1 (2 - x^2)$$  \hspace{1cm} (4.20)

$$+ 2987.83 L_0 + 273.05 L_0^2 - 923.48 L_0^3 - 236.76 L_0^4 - 33.86 L_0^5 - 4 L_0^6$$

and

$$P^{(3)s}_{2}(x) = 19.70002 x_1 L_1 - 3.435474 x_1 L_1^2 + 250(x_1(-4.7656 + 1.6908 x + 0.1703 x^2)$$

$$- 0.41652 x L_0 + 0.90777 x L_0^2 + 0.12478 x L_0^3 + 0.17155 x L_1 + 0.17191 L_0 L_1)$$

$$- 647.3971 L_0 - 66.41219 L_0^2 - 5.353347 L_0^3 - 5.925926 L_0^4 - 0.3950617 L_0^5. \caption{4.21}

The last equation is a high-accuracy parametrization, constructed in the same manner as eq. (4.11) above, of the exact result given in eq. (4.11) of ref. [32].

The trial functions considered for all three cases lead to very similar predictions for the respective next moments, i.e., $N = 18$ for $P^{(3)+}_{\text{ns}}(x)$, and $N = 17 / 19$ for $P^{(3)−/s}_{\text{ns}}(x)$. The residual uncertainty at these $N$-values is a consequence of the width of the bands at large $x$, which in turn (recall figures 5 and 6) is correlated with the uncertainties at smaller $x$. If the spread of the result $A$ and $B$ would underestimate the true remaining uncertainties, then a comparison with additional analytic results at these next values of $N$ should reveal a discrepancy.

In order to check this, we have extended the diagram computations of the $n_f^1$ parts of $P^{(3)−}_{\text{ns}}(x)$ and $P^{(3)+}_{\text{ns}}(x)$ to $N = 17$ and $N = 18$, respectively. The comparison of these results with the Mellin-transformed $n_f^1$ contributions to eqs. (4.12), (4.13), (4.15) and (4.16) yields

$$P^{(3)−}_{\text{ns}}(N = 17) : 194.7126372 B < 194.7126913_{\text{exact}} < 194.7127561_A,$$

$$P^{(3)+}_{\text{ns}}(N = 18) : 195.8888792 B < 195.8888857_{\text{exact}} < 195.888968_A. \caption{4.22}$$

Similar successful checks of our approximation procedure have been carried out for the $n_f^0$ parts of $P^{(3)+}_{\text{ns}}(x)$ and the $n_f^1$ part of $P^{(3)s}_{\text{ns}}(x)$ by deriving less accurate approximations.
using one fewer moment and comparing the results to the now unused highest calculated
moments. As far as we can see from this and other checks, our approximation procedure,
which is of course not mathematically rigorous, does not underestimate the remaining
uncertainties.

5 Numerical implications

We are now ready to address two important applications of our new fourth-order results.
First, as already mentioned above, the large-$x$/large-$N$ limits of the splitting functions
include coefficients that are relevant beyond the evolution of parton distributions: the
(light-like) four-loop cusp anomalous dimension $A_4$ and the $\delta(1-x)$ coefficient $B_4$ for quark
fields. We are now able to provide approximate if rather accurate numerical results for these
coefficients. The obvious second application is a (further) improvement of the perturbative
stability of the evolution of the non-singlet quark distributions over a wide range in $x$.

The analytic large-$n_c$ expression for $A_4$ has been presented in eq. (3.8) above. To-
gether with the approximate results in eqs. (4.12) and (4.13) and the known $n_f^2$ and $n_f^3$
contributions, this yields

$$ A_4 = 20702(2) - 5171.9(2) n_f + 195.5772 n_f^2 + 3.272344 n_f^3 $$

in QCD with $n_f$ quark flavours. The numbers in brackets represent the uncertainty of the
preceeding digit, for which we have increased the spread due to eqs. (4.12) and (4.13) by a
factor of 2. eq. (5.1) leads to

$$ A_4 = 7035(2), 3353(2), 141(2) \text{ for } n_f = 3, 4, 5. $$

For comparison: the corresponding $[1/1]$ Padé approximants used so far are 7849, 4313 and
1553 [98]. The agreement of the actual results with these approximants would be (much)
better without the contributions of the quartic group invariant (see below). A similar situ-
ation has been observed for the four-loop beta function in ref. [109]. The expansion of the
cusp anomalous dimension, now to the fourth order in $\alpha_s$, is given by the very benign series

$$ A_q(\alpha_s, n_f=3) = 0.42441 \alpha_s (1 + 0.72657 \alpha_s + 0.73405 \alpha_s^2 + 0.6647(2) \alpha_s^3 + \ldots), $$
$$ A_q(\alpha_s, n_f=4) = 0.42441 \alpha_s (1 + 0.63815 \alpha_s + 0.50998 \alpha_s^2 + 0.3168(2) \alpha_s^3 + \ldots), $$
$$ A_q(\alpha_s, n_f=5) = 0.42441 \alpha_s (1 + 0.54973 \alpha_s + 0.28403 \alpha_s^2 + 0.0133(2) \alpha_s^3 + \ldots). $$

The corresponding results for the four-loop coefficient $B_4$ in eqs. (2.17) and (4.9) read

$$ B_4 = 23393(10) - 5551(1) n_f + 193.8554 n_f^2 + 3.014982 n_f^3 $$

and

$$ B_q(\alpha_s, n_f=3) = 0.31831 \alpha_s (1 + 0.99712 \alpha_s + 1.24116 \alpha_s^2 + 1.0791(13) \alpha_s^3 + \ldots), $$
$$ B_q(\alpha_s, n_f=4) = 0.31831 \alpha_s (1 + 0.87192 \alpha_s + 0.97833 \alpha_s^2 + 0.5649(13) \alpha_s^3 + \ldots), $$
$$ B_q(\alpha_s, n_f=5) = 0.31831 \alpha_s (1 + 0.74672 \alpha_s + 0.71907 \alpha_s^2 + 0.1085(13) \alpha_s^3 + \ldots). $$
| colour factor | $A_4$ | $B_4$ |
|---------------|-------|-------|
| $C_F^4$       | 0     | 197. ± 3. |
| $C_F^2 C_A$   | 0     | -687. ± 10. |
| $C_F^2 C_A^2$ | 0     | 1219. ± 12. |
| $C_F C_A^3$   | 610.3 ± 0.3 | 295.6 ± 2.4 |
| $d_F^{abcd}d_A^{abcd}/N_R$ | -507.5 ± 6.0 | -996. ± 45. |
| $n_f C_F^3$   | -31.00 ± 0.4 | 81.4 ± 2.2 |
| $n_f C_F^2 C_A$ | 38.75 ± 0.2 | -455.7 ± 1.1 |
| $n_f C_F C_A^2$ | -440.65 ± 0.2 | -274.4 ± 1.1 |
| $n_f d_F^{abcd}d_A^{abcd}/N_R$ | -123.90 ± 0.2 | -143.5 ± 1.2 |
| $n_f^2 C_F^2$ | -21.31439 | -5.775288 |
| $n_f^2 C_F C_A$ | 58.36737 | 51.03056 |
| $n_f^3 C_F$   | 2.454258 | 2.261237 |

**Table 2.** Numerical results for the large-$x$ coefficients $A_4$ and $B_4$ for the seven colour factors contributing to the $n_f^0$ and $n_f$ parts. For completeness also the exactly known $n_f^2$ and $n_f^3$ coefficients are included.

The dominant errors in eq. (5.1) and (5.4) are those of the $n_f$-independent part; its relative uncertainty is $10^{-4}$ for $A_4$ and about four times larger for $B_4$. Due to constraints by large-$N$ moments, the errors of $A_4$ and $B_4$ are fully correlated. The relative uncertainties are larger for the physically relevant values of $n_f$, yet the accuracy in eqs. (5.3) and (5.5) should be amply sufficient for phenomenological applications.

It may be interesting, for theoretical purposes, to consider the contributions of the individual colour factors to $A_4$ and $B_4$. By repeating the approximation procedure of the previous sections separately for each colour factor, we arrive at the corresponding results collected in table 2. Our results show that both quartic group invariants definitely contribute to the four-loop cusp anomalous dimension — an issue that has attracted some interest, see, e.g., refs. [110–114] — which means that the so-called Casimir scaling between the quark and gluon cusp anomalous dimensions, $A_q = C_F/C_A A_g$, does not hold beyond three loops. A lower value, -113.66 after conversion to our notation, results from assumptions made in ref. [115] for the coefficient of $n_f d_F^{abcd}d_A^{abcd}/N_R$.

We now turn to the effect of the four-loop splitting functions (4.6)–(4.21) on the evolution — specifically the logarithmic derivatives $\hat{q}_{\text{ns}}^i \equiv d \ln q_{\text{ns}}^i / d \ln \mu_f^2$ where $\mu_f$ is the factorization scale — of the non-singlet combinations $\hat{q}_{\text{ns}}^{\pm,\nu}(x, \mu_f^2)$ of the quark and anti-quark distributions. In all three cases we employ the same schematic, but characteristic model distribution

$$xq_{\text{ns}}^{\pm,\nu}(x, \mu_f^2) = x^{0.5}(1-x)^3.$$  \hspace{1cm} (5.6)

This facilitates a direct comparison of effects of the various contributions of the splitting
functions. For the same reason the reference scale is specified by the order-independent value

$$\alpha_s(\mu_0^2) = 0.2$$ \hspace{1cm} (5.7)

for the strong coupling constant. This value corresponds to $\mu_0^2 \simeq 25 \ldots 50 \text{GeV}^2$ for $\alpha_s(M_Z^2) = 0.114 \ldots 0.120$ beyond the leading order. In this region of the physical scale $Q^2$ deep-inelastic scattering has been measured both at fixed-target experiments and, for much smaller $x$, at the $ep$ collider HERA. Our default for the number of effectively massless flavours is $n_f = 4$.

The reliability of perturbative calculations can be assessed by the relative size of the higher-order correction at a ‘nominal’ value of the renormalization scale $\mu_r$, here $\mu_r = \mu_f$, and by investigating the stability of the results under variations of $\mu_r$. For $\mu_r \neq \mu_f$ the inverse Mellin transform, see eq. (2.8), of the perturbative expansion (2.9) in terms of $a_s = \alpha_s/(4\pi)$ has to be replaced by

$$P_{\text{ns}}(\mu_f, \mu_r) = a_s(\mu_r^2) P_{\text{ns}}^{(0)} + a_s^2(\mu_r^2) \left( P_{\text{ns}}^{(1),i} - \beta_0 L P_{\text{ns}}^{(0)} \right)$$

$$+ a_s^3(\mu_r^2) \left( P_{\text{ns}}^{(2),i} - 2\beta_0 L P_{\text{ns}}^{(1),i} - \left\{ \beta_1 L - \beta_0^2 L^2 \right\} P_{\text{ns}}^{(0)} \right)$$

$$+ a_s^4(\mu_r^2) \left( P_{\text{ns}}^{(3),i} - 3\beta_0 L P_{\text{ns}}^{(2),i} - \left\{ 2\beta_1 L - 3\beta_0^2 L^2 \right\} P_{\text{ns}}^{(1),i} \right.$$

$$- \left\{ \beta_2 L - \frac{5}{2} \beta_1 \beta_0 L^2 + \beta_0^3 L^3 \right\} P_{\text{ns}}^{(0)} \right) \text{ with } L = \ln \frac{\mu_f^2}{\mu_r^2}. \hspace{1cm} (5.8)$$

For the $\overline{\text{MS}}$ expansion coefficients $\beta_k$ of the beta function of QCD to $N^3\text{LO}$ see refs. [109, 116] and references therein.

In figure 9 the consequences of varying $\mu_r$ over the range $\frac{1}{8} \mu_f^2 \leq \mu_r^2 \leq 8 \mu_f^2$ are displayed for $\dot{q}_{\text{ns}}^i$ at six representative values of $x$ ranging from $x = 0.8$ to $x = 10^{-4}$. A clear improvement of the scale stability to $N^3\text{LO}$ is found over this whole range. Due to the small size of the four-loop contributions and the ‘$x$-averaging’ effect of the Mellin convolution given by

$$[ P_{\text{ns}}^{(n)} \otimes q_{\text{ns}}](x) = \int_x^1 \frac{dy}{y} P_{\text{ns}}^{(n)}(y) q_{\text{ns}} \left( \frac{x}{y} \right) $$ \hspace{1cm} (5.9)

and its generalization for the plus-distribution contributions, the approximate results of section 4 are applicable to lower values of $x$ than one might expect from figure 8.

The relative scale uncertainties of the $\mu_r$-average results, conventionally estimated by

$$\Delta \dot{q}_{\text{ns}}^i = \frac{\max \left[ \dot{q}_{\text{ns}}^i(x, \mu_r^2 = \frac{1}{4} \mu_f^2 \ldots 4 \mu_f^2) \right] - \min \left[ \dot{q}_{\text{ns}}^i(x, \mu_r^2 = \frac{1}{4} \mu_f^2 \ldots 4 \mu_f^2) \right]}{2 \left\{ \text{average} \left[ \dot{q}_{\text{ns}}^i(x, \mu_r^2 = \frac{1}{4} \mu_f^2 \ldots 4 \mu_f^2) \right] \right\} \right| \hspace{1cm} (5.10)$$

is shown in the left panels of figures 10, 11 and 12 for all three cases $i = +$, $-$ and $v$. In the corresponding right panels, the relative size of the $N^3\text{LO}$ corrections at the scale $\mu_r = \mu_f$ are compared to the relative $N^2\text{LO}$ effects. Both the relative scale uncertainties and the relative corrections have a singularity at about $x \simeq 0.07$ due to a sign change of the scaling violations $d\dot{q}_{\text{ns}}^i(x)/d\ln \mu_f^2$. 

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Outside the region around $x = 0.07$ where the $\mu^-$-derivatives are small, the remaining uncertainty of $q_{ns}^+$ is well below 1% down to $x \simeq 10^{-3}$ and possibly, below. The size and $\mu^-$-variation of the NLO and NNLO contributions are somewhat larger for $q_{ns}^-$ at small $x$, yet neither the N$^3$LO correction nor its scale variation exceeds 1% in the region shown in the plot.

The case of $q_{ns}^-$, shown in figure 12 is noticeably different beyond NLO due to the appearance of the $d^{abc}d_{abc}\eta_f$ contribution $P_{ns}^\pm$ which is negligible and at large $x$ but large at small $x$ [14]: the difference of the NNLO curves of figure 11 and 12 — note the different scales for the ordinate — is caused completely by this contribution due to our choice (5.6) of the input quark distribution. Also in this case our new N$^3$LO results leads to a considerable improvement and a remaining small-$x$ uncertainty of no more than about 2% at $x \geq 10^{-4}$.

6 The time-like case

The differences between the initial-state (‘space-like’, $\sigma = -1$) and the final-state (‘time-like’, $\sigma = 1$) splitting functions respectively governing the evolution of the parton distributions and fragmentation functions can be expressed in terms of lower-order quantities. At N$^3$LO they read
\[
\delta P^{(3)}(x) \equiv P_{ns,\sigma=1}^{(3)}(x) - P_{ns,\sigma=-1}^{(3)}(x) = \\
= 2 \left\{ [\ln x \cdot P_{ns}^{(2)}] \otimes P_{ns}^{(0)} + [\ln x \cdot P_{ns}^{(0)}] \otimes P_{ns}^{(2)} + [\ln x \cdot P_{ns}^{(1)}] \otimes P_{ns}^{(1)} \right\} \\
- 2 P_{ns}^{(0)} \otimes [\ln x \cdot P_{ns}^{(0)}] \otimes 3 - 4 [P_{ns}^{(0)}] \otimes 3 \otimes [\ln x \cdot P_{ns}^{(0)}] \otimes [\ln^2 x \cdot P_{ns}^{(0)}] \\
- \frac{2}{3} [P_{ns}^{(0)}] \otimes 3 \otimes [\ln^3 x \cdot P_{ns}^{(0)}],
\]
where we have used the short-hand notations $A \otimes A \equiv A \otimes A$ etc for the Mellin convolutions, and $P_{ns}^{(n)}$ stands for the average of the corresponding $\sigma = 1$ and $\sigma = -1$ expansion coefficients, normalized as in eq. (2.9). Eq. (6.1) has been derived in ref. [56] by generalizing results in ref. [65]; it is also a direct consequence of eq. (2.15) [67].

The resulting rather lengthy explicit expressions can be found in appendix C. Here we present parametrizations in terms of powers of $x$ and the logarithms in eq. (4.10). As above, their small-$x$ and large-$x$ coefficients are exact up to their rounding to seven digits. The accuracy of the $n_f$ coefficients is better than 0.1% except close to zeros. The three parametrizations are given by
\[
\delta P^{(3)+}(x) = \\
-33901.87 L_1 - 32392.47 + 25000(x(1.2960 + 1.7438x - 1.0943x^2 - 0.44064x^3)) \\
+ x L_0(0.6440 + 0.8939 L_0 + 0.21405 L_0^2) + L_1 (2.0343 x_1 + 0.35738 L_0) - 10399.74 L_0 \\
- 25718.24 L_0^3 - 5965.487 L_0^3 - 206.7846 L_0^4 + 4.213992 L_0^5 - 0.7023320 L_0^6 \\
+ n_f \left[ 5483.660 L_1 + 4975.255 + 250(x(-19.877 - 8.0977 x + 12.335 x^2 + 8.1174 x^3)) \\
+ x L_0 (13.617 - 7.8856 L_0 - 2.2491 L_0^2) + L_1 (-20.171 x_1 + 6.571 L_0) + 657.5425 L_0
\]
Figure 9. The dependence of the NLO, NNLO and N$^3$LO predictions for $\dot{q}_{\text{NS}}^+ \equiv d \ln q_{\text{NS}}^+/d \ln \mu_f^2$ on the renormalization scale $\mu_r$ at six typical values of $x$ for the initial conditions (5.6) and (5.7). The effect of the remaining uncertainty of the four-loop splitting function $P_{\text{NS}}^{(3)+}$, indicated by the difference of the solid and dotted curves, is practically invisible except for the last two panel.
\begin{align*}
\Delta \dot{q}_{\text{NS}}^+ = & \quad -33901.87L_1 - 32392.47 + 25000(x_1(1.2892 + 1.4892x - 1.4262x^2 + 0.29374x^3)) \\
+ & \quad xL_0(1.1307 + 0.17484L_0 + 0.14894L_0^2 + L_1(5.0547x_1 + 3.4824L_0)) - 7307.364L_0 \\
- & \quad 24617.82L_0^2 - 7051.323L_0^3 - 665.0339L_0^4 - 13.82716L_0^5 + 1.035940L_0^6 \\
+ n_f^2 \left( & -53.37723L_1 - 80.32433 + 50(x_1(1.6030 + 15.938x - 5.3145x^2 + 1.8682x^3)) \\
+ & \quad xL_0(13.301 + 2.1060L_0 + 0.4375L_0^2 + L_1(13.060x_1 + 11.023L_0)) - 9.550559L_0 \\
- & \quad 56.98805L_0^2 - 21.59671L_0^3 - 1.580247L_0^4 \right), 
\end{align*}

\begin{equation}
\delta P^{(3)}(x) = 
-33901.87L_1 - 32392.47 + 25000(x_1(1.2892 + 1.4892x - 1.4262x^2 + 0.29374x^3)) \\
+ xL_0(1.1307 + 0.17484L_0 + 0.14894L_0^2 + L_1(5.0547x_1 + 3.4824L_0)) - 7307.364L_0 \\
- 24617.82L_0^2 - 7051.323L_0^3 - 665.0339L_0^4 - 13.82716L_0^5 + 1.035940L_0^6 \\
+ n_f^2 \left( & -53.37723L_1 - 80.32433 + 50(x_1(1.6030 + 15.938x - 5.3145x^2 + 1.8682x^3)) \\
+ & \quad xL_0(13.301 + 2.1060L_0 + 0.4375L_0^2 + L_1(13.060x_1 + 11.023L_0)) - 9.550559L_0 \\
- & \quad 56.98805L_0^2 - 21.59671L_0^3 - 1.580247L_0^4 \right)
\end{equation}
The difference of the time-like and space-like splitting functions and the resulting scale derivative are illustrated in figures 13 and 14 for the most important case, NS\(^+\). The pattern is somewhat different in the time-like case, e.g., the N\(^3\)LO contribution is negative at small \(x\). Yet also here the perturbative expansion is ‘perfectly’ stable after including the N\(^3\)LO corrections, with a residual uncertainty of 1% or less for \(x > 10^{-4}\) at our (for the time-like case, low-scale) reference point.
Figure 12. As figures 10 and 11, but for the total (flavour-summed) valence quark distribution $q_{NS}^v$. Here and in figures 10 and 11 the spikes close to $x = 0.07$ reflect the sign-changes of $\dot{q}_{NS}$ and do not constitute appreciable absolute corrections and uncertainties. All three figures refer to $n_f = 4$ and $\alpha_s(\mu_r^2 = \mu_f^2) = 0.2$.

7 Summary and outlook

We have presented the four-loop corrections $P_{ns}^{(3)\pm,v}(x)$ to all three non-singlet splitting functions in perturbative QCD. Our results, which are partly approximate but sufficiently accurate for collider-physics applications, allow to set-up and solve the QCD evolution equations for flavour non-singlet (initial-state) parton distributions (PDFs) and (final-state) fragmentation functions (FFs) at $N^3$LO. They thus provide a major step towards the consistent application of QCD factorization to theoretical predictions for $N^3$LO cross sections with initial (final) state hadrons, as already obtained in refs. [21–24, 26], which requires hard partonic cross section and PDFs (FFs) at the same order in renormalization-group improved perturbation theory.

The resulting logarithmic scale derivatives $d\ln q_{ns}^{\pm,v}/d\ln \mu_r^2$ exhibit a very good convergence of the perturbative expansion. Both the four-loop corrections and the $N^3$LO dependence on the renormalization scale $\mu_r$ mostly amount to as little as 1% or less (and maximally 2%, for $P_{ns}^v(x)$ at small $x$) at momentum fractions $x > 10^{-4}$ for $\alpha_s(\mu_f^2) \simeq 0.2$.

Our results have been obtained via computations of fixed Mellin moments — to $N = 20$ for the diagrams contributing in the limit of a large number $n_c$ of colours, and $N = 16$ otherwise — for the QCD corrections to quark and anti-quark operator matrix elements (OMEs) up to four loops. After projection of all external spins and Lorentz indices, these
Figure 13. The perturbative expansion of the difference \( \delta P^+ = P_{\text{ns}, \sigma=1}^+ - P_{\text{ns}, \sigma=-1}^+ \) of the time-like (\( \sigma = 1 \)) and space-like (\( \sigma = -1 \)) non-singlet splitting functions. Left: results in Mellin-space, compared to beyond-LO contributions in the space-like case. Right: convolution with the schematic input shape (5.6).

OMEs reduce to four-loop massless propagator integrals which can be evaluated with the \textsc{Forcer} program \cite{33} in the computer algebra system \textsc{Form} \cite{34,35,74}.

For the large-\( n_c \) contributions to \( P_{\text{ns}}^{(3)\pm} \), these moments turn out to be sufficient for a reconstruction of the all-\( N \) results in terms of harmonic sums by solving systems of Diophantine equations. Additional knowledge about the limits for \( x \to 0 \) and \( x \to 1 \), as well as the rephrasing (based on conformal symmetry) of the evolution equations in terms of a universal ‘reciprocity-respecting’ evolution kernel \cite{65,66,67}, have been instrumental in this step. Beyond the large-\( n_c \) contributions we have used the computed Mellin moments, again supplemented by endpoint constraints, to provide approximations for the four-loop splitting functions including \( x \)-dependent estimates of their residual uncertainties. The latter are very small, except in the region \( x < 10^{-2} \). These small-\( x \) uncertainties are subject to a further suppression in the actual evolution due to the convolution with the PDFs (or FFs). Due to this our results are found to be sufficiently precise for \( x \gtrsim 10^{-4} \).

From the threshold limit \( x \to 1 \) we have been able to determine the complete cusp anomalous dimension \( A_4 \) for quarks at four loops. Our results for the \( n_f^0 \) and \( n_f^1 \) parts beyond the large-\( n_c \) contribution are numerical, and lead to a relative accuracy of \( 10^{-4} \) for these coefficients in QCD, which should be amply sufficient for phenomenological applications. The break-up of \( A_4 \) in terms of individual colour factors includes non-vanishing contributions with the quartic group invariants. The exact results for \( A_4 \) of the present
Figure 14. As figure 10, but for its time-like counterpart, the fragmentation function $f^{+}_{\text{NS}}$. The same input (5.6) and (5.7) is used here, so all differences to figure 10 are due to the different $N^{n>0}\text{LO}$ splitting functions.

The terms $B_4$ proportional to $\delta(1-x)$ in the four-loop splitting functions yield the universal eikonal anomalous dimension if properly combined with information on infrared singularities from the QCD form factor. This is an important ingredient for extending the threshold resummation for inclusive cross section to $N^4\text{LL}$ accuracy, i.e. (next-to-)4-leading logarithmic order.

In order to practically complete the QCD evolution equations at $N^3\text{LO}$, corresponding results are required for the singlet splitting functions at four loops, i.e., the pure-singlet quark-quark splitting function $P^{(3)}_{\text{ps}}(x)$ and those involving gluons, $P^{(3)}_{\text{qg}}(x)$, $P^{(3)}_{\text{gq}}(x)$ and $P^{(3)}_{\text{gg}}(x)$. All these quantities are currently unknown beyond the $N = 2$ and $N = 4$ moments presented in ref. [31]. However, by following the approach of the present paper, it should be feasible to compute enough moments of the (theoretically much more complicated, see refs. [12, 119]) corresponding flavour-singlet OMEs up to four loops with the FORCER program to gather sufficient information for first phenomenologically relevant approximations. We leave this topic to future research.
Incremental improvements in the flavour non-singlet sector can be obtained by calculating more moments, which is a hard problem within the present computational set-up for almost all colour factors, and by incorporating more external information, such as, e.g., a future exact result for the four-loop cusp anomalous dimension from calculations of the photon-quark form factor.

A derivation of the exact expressions for the $n_f$-independent hardest parts will require, in addition, a much improved theoretical understanding. In this context it may be interesting to note that the $\zeta_5$ part of $\gamma_{\text{ns}}^{(3)}(N)$, which can be determined at all $N$ from the presently available information (see appendix D) includes a contribution

$$-\frac{128}{3} \left\{ 3 C_F^2 C_A^2 - 2 C_F C_A^3 + 12 \frac{d_F^{abcd} d_F^{abcd}}{N_R} \right\} 5 \zeta_5 [S_1(N)]^2 \quad \text{(7.1)}$$

that vanishes in the large-$n_c$ limit. The resulting $\ln^2 N$ large-$N$ behaviour needs to be compensated by non-$\zeta_5$ terms, and it is tempting to identify the $5 \zeta_5$ in eq. (7.1) as the $\zeta_5$-'tail’ of the function

$$f(N) = 5 \zeta_5 - 2 S_{-5} + 4 S_{-2} \zeta_3 - 4 S_{-2,-3} + 8 S_{-2,-2,1} + 4 S_{3,-2} - 4 S_{4,1} + 2 S_5 \quad \text{(7.2)}$$

This function first occurred multiplied with positive powers of $N$ in three-loop coefficient functions of DIS in ref. [22] and resurfaced, now multiplied with $[S_1(N)]^2$ as in eq. (7.1), as the ‘wrapping correction’ in the anomalous dimensions in $\mathcal{N} = 4$ maximally supersymmetric Yang-Mills theory [120], where it is crucial for obtaining the correct small-$x$ limit, see, e.g., ref. [121]. We thus hypothesize that eq. (7.1) represents the first glimpse of the wrapping corrections in an anomalous dimension in QCD.

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A  Feynman rules

Below we present the Feynman rules for vertices arising from insertions of the operator $O_{\{\mu_1,..,\mu_N\}}^n$ in eq. (2.1). All momenta $p_i$ are flowing into the operator vertex and we use

$$q = \sum_i p_i ,$$

(A.1)

where $q$ is the outgoing momentum flow through the operator. The free Lorentz indices of the operator $O_{\{\mu_1,..,\mu_N\}}^n$ are contracted with

$$\Delta_{\mu_1} \cdots \Delta_{\mu_N} ,$$

where the vector $\Delta$ fulfils $\Delta^2 = 0$. We limit the derivation up to four additional gluons coupling to the operator, i.e., $n = 6$ in figure 1. For Feynman rules with up to three additional gluons and zero momentum flow through the operator, see also ref. [72] and references therein.

The expressions for unpolarized quark operators in eqs. (A.2)–(A.6) are readily generalized to the polarized case by substituting $\Delta \rightarrow \Delta \gamma_5$.

\[
\begin{align*}
\Delta (\Delta \cdot p_2)^{N-1} \\
\end{align*}
\] (A.2)

\[
\begin{align*}
-g t^{a_3} \Delta \Delta^{\mu_3} \sum_{j_1=0}^{N-2} (p_2 \cdot \Delta)^{N-2-j_1} (q \cdot \Delta - p_1 \cdot \Delta)^{j_1} \\
\end{align*}
\] (A.3)

\[
\begin{align*}
g^2 \Delta \Delta^{\mu_3} \Delta^{\mu_4} \sum_{j_1=0}^{N-3} \sum_{j_2=0}^{j_1} (p_2 \cdot \Delta)^{N-3-j_1} (q \cdot \Delta - p_1 \cdot \Delta)^{j_1} \left( t^{a_3} t^{a_4} (q \cdot \Delta - p_1 \cdot \Delta - p_3 \cdot \Delta)^{j_1-j_2} \right) \\
\end{align*}
\]
\[ t^{a_4 t^{a_3}} (q \cdot \Delta - p_{1} \cdot \Delta - p_{4} \cdot \Delta)^{j_1 - j_2} \]  
(A.4)

\[ - g^{3} \Delta^{\mu_3} \Delta_{\mu_4} \Delta_{\mu_5} \sum_{j_1=0}^{N-4} \sum_{j_2=0}^{j_1} \sum_{j_3=0}^{j_2} (p_2 \cdot \Delta)^{N-4-j_1} (q \cdot \Delta - p_{1} \cdot \Delta)^{j_3} \]

\[ \left( t^{a_3 t^{a_4} t^{a_5} t^{a_6}} (q \cdot \Delta - p_{1} \cdot \Delta - p_{3} \cdot \Delta)^{j_2 - j_3} (q \cdot \Delta - p_{1} \cdot \Delta - p_{4} \cdot \Delta)^{j_1 - j_2} + t^{a_3 t^{a_4} t^{a_5} t^{a_6}} (q \cdot \Delta - p_{1} \cdot \Delta - p_{4} \cdot \Delta - p_{5} \cdot \Delta)^{j_1 - j_2} + t^{a_4 t^{a_5} t^{a_6}} (q \cdot \Delta - p_{1} \cdot \Delta - p_{5} \cdot \Delta)^{j_2 - j_3} + t^{a_3 t^{a_4} t^{a_5} t^{a_6}} (q \cdot \Delta - p_{1} \cdot \Delta - p_{5} \cdot \Delta - p_{6} \cdot \Delta)^{j_1 - j_2} \right) \quad \text{(A.5)} \]
\[ (q \cdot \Delta - p_1 \cdot \Delta - p_3 \cdot \Delta - p_4 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \\
+ l_4^{a_3} l_6^{a_5} l_4^{a_5} l_6^{a_4} (q \cdot \Delta - p_1 \cdot \Delta - p_3 \cdot \Delta)^{j_1-j_2} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_2-j_3} \\
+ l_4^{a_4} l_6^{a_3} l_4^{a_5} l_6^{a_6} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \\
+ l_4^{a_4} l_6^{a_6} l_4^{a_5} l_6^{a_3} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \\
+ l_4^{a_4} l_6^{a_4} l_4^{a_6} l_6^{a_3} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \\
+ l_4^{a_4} l_6^{a_4} l_4^{a_6} l_6^{a_5} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \\
+ l_4^{a_5} l_6^{a_4} l_4^{a_6} l_6^{a_3} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \\
+ l_4^{a_5} l_6^{a_4} l_4^{a_6} l_6^{a_5} (q \cdot \Delta - p_1 \cdot \Delta - p_4 \cdot \Delta - p_5 \cdot \Delta - p_6 \cdot \Delta)^{j_1-j_2} \]
Here we present the anomalous dimensions at four loops for $1 \leq N \leq 16$. Obviously, $\gamma_{\bar{m}s}(N = 1) = 0$ at all orders. To fix our normalization, we write down the complete expression for $N = 2$ including all lower orders. Recall that $a_s = \alpha_s/(4\pi)$.

\begin{align}
\gamma_{\bar{m}s}(N=2) &= a_s \left\{ \frac{8}{3} C_F + a_s^2 \left\{ \frac{112}{27} C_F^2 - \frac{376}{27} C_A C_F - \frac{64}{27} n_f C_F \right\} \right.
onumber \\
&\quad + a_s^3 \left\{ C_F^3 - \frac{560}{243} + \frac{128}{3} \zeta_3 \right\} + C_A C_F^2 \left[ \frac{8528}{243} - 64 \zeta_3 \right] + C_A^2 C_F \left[ \frac{20920}{243} + 64 \frac{4}{3} \zeta_3 \right] \\
&\quad + n_f C_F^3 \left[ \frac{3412}{243} + 64 \frac{4}{3} \zeta_3 \right] + n_f C_A C_F \left[ \frac{3128}{243} - \frac{64}{3} \zeta_3 \right] + n_f^2 C_F \left[ \frac{224}{243} \right] \right. \\
&\quad \left. + a_s^4 \left\{ C_A^4 \left[ \frac{194392}{2187} + \frac{10880}{81} \zeta_3 - \frac{1280}{3} \zeta_5 \right] + C_A C_F \left[ \frac{238676}{2187} + \frac{31040}{81} \zeta_3 - \frac{704}{3} \zeta_4 + \frac{1280}{3} \zeta_5 \right] \right. \\
&\quad \left. + C_A^2 C_F^2 \left[ \frac{1626064}{2187} - \frac{25744}{27} \zeta_3 + 352 \zeta_4 + \frac{4480}{9} \zeta_5 \right] \right. \\
&\quad \left. + C_A^3 C_F \left[ \frac{1734130}{2187} + \frac{34936}{81} \zeta_3 - \frac{352}{3} \zeta_4 - \frac{12160}{27} \zeta_5 \right] \right. \\
&\quad \left. + n_f C_F^3 \left[ \frac{190912}{2187} - \frac{2528}{81} \zeta_3 + \frac{128}{3} \zeta_4 - \frac{640}{3} \zeta_5 \right] \right. \\
&\quad \left. + n_f C_A C_F^2 \left[ \frac{177748}{2187} + \frac{12928}{27} \zeta_3 - \frac{544}{3} \zeta_4 + \frac{320}{9} \zeta_5 \right] \right. \\
&\quad \left. + n_f^2 C_A C_F \left[ \frac{53018}{243} - \frac{4040}{9} \zeta_3 + \frac{416}{3} \zeta_4 + \frac{4480}{27} \zeta_5 \right] + n_f^2 C_F^2 \left[ \frac{24944}{2187} - \frac{128}{3} \zeta_3 + \frac{64}{3} \zeta_4 \right] \right. \\
&\quad \left. + n_f^3 C_F \left[ \frac{1024}{2187} + \frac{128}{81} \zeta_3 \right] \right. \\
&\quad \left. + a_F^{abcd} a_F^{abcd} \frac{9}{N_R} \left[ \frac{736}{9} + \frac{1984}{9} \zeta_3 + \frac{5120}{9} \zeta_5 \right] \right. \\
&\quad \left. + n_f a_F^{abcd} a_F^{abcd} \frac{9}{N_R} \left[ \frac{832}{9} + \frac{1024}{9} \zeta_3 - \frac{2560}{9} \zeta_5 \right] \right. \\
&\quad \left. \right\}.
\end{align}
The four-loop coefficient for the even moments $N = 4$ to $N = 16$ are given by

$$\gamma_{ns}^{(3)}(N=4) =
C_F^2 \left[ \frac{3482407012657}{349920000000} + \frac{14504764}{50625} \zeta_3 - \frac{25136}{45} \zeta_5 \right] +
C_A C_F^2 \left[ \frac{33802068299}{174960000} - \frac{10215349}{81000} \zeta_3 - \frac{15829}{75} \zeta_4 + \frac{33004}{45} \zeta_5 \right] +
C_A^2 C_F^2 \left[ -\frac{1557367902137}{174960000} - \frac{2497339}{16875} \zeta_3 + \frac{15829}{50} \zeta_4 - \frac{1645}{9} \zeta_5 \right] +
C_A^3 C_F \left[ \frac{49455970561}{43740000} + \frac{13461191}{81000} \zeta_3 - \frac{15829}{150} \zeta_4 - \frac{18646}{135} \zeta_5 \right] +
n_f C_F^2 \left[ \frac{29581840417}{174960000} + \frac{820961}{10125} \zeta_3 + \frac{2878}{75} \zeta_4 - \frac{1256}{3} \zeta_5 \right] +
n_f C_A C_F^2 \left[ -\frac{89325051233}{437400000} + \frac{4588639}{6750} \zeta_3 - \frac{21587}{75} \zeta_4 + \frac{628}{9} \zeta_5 \right] +
n_f C_A^2 C_F \left[ -\frac{1796654459}{4860000} - \frac{5247961}{6750} \zeta_3 + \frac{18709}{75} \zeta_4 + \frac{43736}{135} \zeta_5 \right] +
n_f^2 C_F^2 \left[ \frac{5419760639}{21870000} - \frac{2146}{25} \zeta_3 + \frac{628}{15} \zeta_4 \right] + n_f^2 C_A C_F \left[ \frac{60167591}{3645000} + \frac{2146}{25} \zeta_3 - \frac{628}{15} \zeta_4 \right] +
n_f^3 C_F \left[ \frac{17813699}{21870000} + \frac{1256}{405} \zeta_3 + \frac{d_{C F}^{d C F}}{N R} \left[ \frac{254713}{1350} + \frac{63568}{45} \zeta_3 - \frac{78868}{45} \zeta_5 \right] \right] +
n_f^4 C_F \left[ \frac{16568}{135} + \frac{22552}{75} \zeta_3 - \frac{26912}{45} \zeta_5 \right], \quad (B.2)$$

$$\gamma_{ns}^{(3)}(N=6) =
C_F^4 \left[ \frac{287417623549488131}{180108854100000} + \frac{33026498018}{121550625} \zeta_3 - \frac{450712}{735} \zeta_5 \right] +
C_A C_F^3 \left[ \frac{9809771626657}{476478450000} - \frac{9025033804}{121550625} \zeta_3 - \frac{768482}{3675} \zeta_4 + \frac{621976}{735} \zeta_5 \right] +
C_A^2 C_F^2 \left[ -\frac{588570119595401}{91892275000} + \frac{3466052671}{27011250} \zeta_3 + \frac{384241}{1225} \zeta_4 - \frac{1695382}{2205} \zeta_5 \right] +
C_A^3 C_F \left[ \frac{64997866579309}{56010528000} + \frac{395253829}{27783000} \zeta_3 - \frac{384241}{3675} \zeta_4 + \frac{1378042}{6615} \zeta_5 \right] +
n_f C_F^3 \left[ \frac{2697261071752787}{1286498150000} + \frac{2830802}{19845} \zeta_3 + \frac{139724}{3675} \zeta_4 - \frac{11344}{21} \zeta_5 \right] +
n_f C_A C_F^2 \left[ -\frac{160989027854717}{612615150000} + \frac{8830621}{11025} \zeta_3 - \frac{1301446}{3675} \zeta_4 + \frac{5672}{63} \zeta_5 \right]$$
\[\gamma_{\text{ns}}^{(3)}(N = 8) = \]
\[
\frac{1526099947627950150545831}{8067032349014106000000} + \frac{20349260276089}{787648050000} \zeta_3 - \frac{1804282}{2835} \zeta_5
\]
\[
+ \frac{142175220996188451677}{19207219878604800000} \zeta_3 - \frac{338728975833}{1400263200000} \zeta_3 + \frac{27614477}{132300} \zeta_4 + \frac{17728727}{19845} \zeta_5
\]
\[
+ \frac{13695011182141605151}{3048765060096000000} \zeta_3 + \frac{8102719994761}{2100394800000} \zeta_3 + \frac{27614477}{88200} \zeta_4 - \frac{4834637}{3780} \zeta_5
\]
\[
+ \frac{100893395789394973}{8710757314560000} - \frac{33568622483}{244944000} \zeta_3 - \frac{27614477}{264600} \zeta_4 + \frac{31576003}{59535} \zeta_5
\]
\[
+ \frac{113630909563809311301}{48018049696512000000} + \frac{5612836699}{30005640} \zeta_4 + \frac{2510407}{66150} \zeta_4 - \frac{39532}{63} \zeta_5
\]
\[
+ \frac{11533221344476533811}{381095632512000000} + \frac{88755158137}{100018800} \zeta_3 - \frac{17730227}{44100} \zeta_4 + \frac{19766}{189} \zeta_5
\]
\[
+ \frac{189224450730372949}{36294822414000000} - \frac{1661823807613}{1500282000} \zeta_4 + \frac{48169867}{813749} \zeta_4 - \frac{132300}{1701} \zeta_5
\]
\[
+ \frac{386628670506434189}{9527390812800000} - \frac{77750543}{595350} \zeta_3 + \frac{19766}{315} \zeta_4
\]
\[
+ \frac{2001327622434827}{907370553600000} + \frac{77750543}{595350} \zeta_3 - \frac{19766}{315} \zeta_4
\]
\[
+ \frac{113630909563809311301}{48018049696512000000} + \frac{5612836699}{30005640} \zeta_4 + \frac{2510407}{66150} \zeta_4 - \frac{39532}{63} \zeta_5
\]
\[
+ \frac{11533221344476533811}{381095632512000000} + \frac{88755158137}{100018800} \zeta_3 - \frac{17730227}{44100} \zeta_4 + \frac{19766}{189} \zeta_5
\]
\[
+ \frac{189224450730372949}{36294822414000000} - \frac{1661823807613}{1500282000} \zeta_4 + \frac{48169867}{813749} \zeta_4 - \frac{132300}{1701} \zeta_5
\]
\[
+ \frac{386628670506434189}{9527390812800000} - \frac{77750543}{595350} \zeta_3 + \frac{19766}{315} \zeta_4
\]
\[ \gamma_{\text{ns}}^{(3)}(N = 10) = \]
\[ C_F \left[ \frac{2017383724760695233171991134991}{9825227427985488199296000000} + \frac{143667693462054187}{576597755002500} \zeta_3 \right. \]
\[ - \frac{222343766}{343035} \zeta_5 \]
\[ + C_A C_F^3 \left[ - \frac{151796299}{727650} \zeta_4 + \frac{440406986}{480249} \zeta_5 \right] \]
\[ + C_A^2 C_F^2 \left[ - \frac{8344774766025843631923173}{278981057780409600000000} + \frac{488048876510787149}{768797006670000} \zeta_3 - \frac{151796299}{1455300} \zeta_4 \right. \]
\[ + \frac{2378996101}{2881494} \zeta_5 \]
\[ + n_f C_F^3 \left[ \frac{2730954981137825155546834719}{10633630173003119040000000} + \frac{110938456177788}{499218835500} \zeta_3 \right. \]
\[ + \frac{151796299}{4002075} \zeta_4 - \frac{482200}{693} \zeta_5 \]
\[ + n_f C_A C_F^2 \left[ - \frac{232521423145307204017937}{69745264445102400000000} + \frac{317697741600923}{332812557000} \zeta_3 \right. \]
\[ - \frac{1172854799}{2668050} \zeta_4 + \frac{241100}{2079} \zeta_5 \]
\[ + n_f C_A^2 C_F \left[ - \frac{3449950935218411700067}{60385510342080000000} + \frac{168584409069107}{138671898750} \zeta_3 \right. \]
\[ + \frac{3214971799}{8004150} \zeta_4 + \frac{181204484}{343035} \zeta_5 \]
\[ + n_f^2 C_F^2 \left[ \frac{17653381355656128339793}{3853989544480632000000} - \frac{5236700507}{36018675} \zeta_3 + \frac{48220}{693} \zeta_4 \right. \]
\[ + n_f^2 C_A C_F \left[ \frac{7133251712396708693}{3019275517104000000} + \frac{5236700507}{36018675} \zeta_3 - \frac{48220}{693} \zeta_4 \right. \]
\[ + n_f^3 C_F \left[ - \frac{613551152411968391}{4981804603221600000} + \frac{96440}{18711} \zeta_3 \right. \]
\[
\begin{align*}
+ \frac{d_Fa_F^3}{N_R} a_F^3 & \left[ \frac{6387465101013091}{41493513600000} + \frac{751245802206203}{15127843500000} \zeta_3 - \frac{18465293144}{2401245} \zeta_5 \right] \\
+ n_f & \frac{d_Fa_F^3}{N_R} a_F^3 \left[ \frac{51618947156153}{134469720000} + \frac{18626595374}{33350625} \zeta_3 - \frac{141624128}{114345} \zeta_5 \right]
\end{align*}
\]
\[\gamma^{(3)}_{n_{\text{ns}}} \quad (N = 12) = \]
\[
C_F^3 \left[ \frac{331516931678819616019516119238120540753}{1514296125734353420506706111008000000} + \frac{3996930795976177310551}{16468208480626402500} \zeta_3 \\
- \frac{39764613926}{57972915} \right] \\
+ C_A C_F^3 \left[ - \frac{4110347217755725651032771751930697}{2566260615608615060517168000000} + \frac{98597799463449331819}{274470141343773750} \zeta_4 \\
- \frac{25648239313}{122972850} \zeta_4 + \frac{34328650826}{36891855} \zeta_5 \right] \\
+ C_A^2 C_F^2 \left[ - \frac{23885109171514866677698033404359}{1367305023301293405288960000000} + \frac{38330693033675619062839}{43915222615003740000} \zeta_3 \\
+ \frac{25648239313}{81981900} \zeta_4 + \frac{27699141079}{12882870} \zeta_5 \right] \\
+ C_A^3 C_F \left[ \frac{516560961693213572859380630401}{45531302807235877632000000} - \frac{3552780661279687493}{8356447859760000} \zeta_3 \\
- \frac{25648239313}{245945700} \zeta_4 + \frac{242831514391}{221351130} \zeta_5 \right] \\
+ n_f C_F^3 \left[ \frac{14007724721120701350168397616773189}{5135212312172301201034336000000} + \frac{275829519782497567}{1096783781593500} \zeta_3 \\
+ \frac{25648239313}{67650675} \zeta_4 - \frac{6774784}{9009} \zeta_5 \right] \\
+ n_f C_A C_F^2 \left[ - \frac{12260374152873668292904144517329}{34182625582532335132224000000} + \frac{246056489027368057}{243729729243000} \zeta_3 \\
- \frac{21241105873}{450900450} \zeta_4 + \frac{3387392}{27027} \zeta_5 \right] \\
+ n_f C_A^2 C_F \left[ - \frac{581173376177155203686777699}{94856880848080784000000} - \frac{119245993023506038}{91398648466125} \zeta_3 \\
+ \frac{585126838993}{1352701350} \zeta_4 + \frac{599861228}{1054053} \zeta_5 \right] \\
+ n_f^2 C_F^2 \left[ \frac{14406272473786523198816959409}{2848552131877694594352000000} - \frac{191866603189}{1217431215} \zeta_3 + \frac{3387392}{45045} \zeta_4 \right]
\]
\begin{align}
+ n_f^2 C_A C_F \left[ \frac{442446968888889719821403}{1778565515907651470000} + \frac{191866603189}{1217431215} \zeta_3 - \frac{3387392}{45045} \zeta_4 \right] + n_f^3 C_F \left[ \frac{1863394741992678675203}{14228532127621176000} + \frac{6774784}{1216215} \zeta_3 \right] \\
+ \frac{d_F^\text{bed} d_F^\text{bed}}{N_R} \left[ \frac{55568956318895924814773}{2915757499374220000} + \frac{224835066938602217}{37496881422000} \zeta_3 - \frac{3785627672012}{405810405} \zeta_3 \right] \\
+ n_f^2 C_A C_F \left[ \frac{22640351375268075599}{4958771257440000} + \frac{18402607883189}{30435780375} \zeta_3 - \frac{485441696}{351351} \zeta_5 \right].
\end{align}

\begin{align}
\gamma_{ns}^{(3)+}(N=14) = \\
C_F^4 \left[ \frac{97523736280058513278817419051773381749}{4403703216384383440387460288000000} + \frac{1567683874499713548181}{658728339225056100} \zeta_3 \right] - \frac{267445675058}{405810405} \zeta_5 \\
+ C_A C_F^3 \left[ \frac{333627771877952197537388891905440907}{17963824309260305420620176000000} + \frac{11882737164573624768177}{2195761130750187000} \zeta_3 \right] - \frac{3663695353}{17567550} \zeta_4 + \frac{6927566278}{7378371} \zeta_5 \\
+ C_A C_F^2 \left[ \frac{86656353251877426313613209043}{128842204118775724729152000000} + \frac{6031813059800336882399}{5489402862875467500} \zeta_3 \right] \\
+ \frac{3663695353}{11711700} \zeta_4 - \frac{45655020935}{18036018} \zeta_5 \\
+ C_A^2 C_F \left[ \frac{3974468919372874040522232093269}{35493720142913422790400000000} - \frac{8803366113150246019}{15746090197240000} \zeta_3 \right] \\
+ \frac{3663695353}{35135100} \zeta_4 + \frac{59678850515}{44270226} \zeta_5 \\
+ n_f C_F^3 \left[ \frac{1469298768736788933360199563898027}{513252132172301201034336000000} + \frac{303399887420914357}{1096783781593500} \zeta_3 \right] \\
+ \frac{3663695353}{96621525} \zeta_4 - \frac{7204840}{9009} \zeta_5 \\
+ n_f C_A C_F \left[ \frac{22747053762883618602535741424929}{59819594769431586481392000000} + \frac{257384193913101469}{243729729243000} \zeta_3 \right] \\
- \frac{31996728653}{64414350} \zeta_4 + \frac{3602420}{27027} \zeta_5 \\
+ n_f C_A^2 C_F \left[ \frac{15060401544121460041592815421}{23239935807589979208000000} - \frac{17663095878365457683}{12795810785257500} \zeta_3 \right] \\
\end{align}
\[\gamma_{\text{ns}}^{(3)}(N = 16) = \]
\[C_{A}^2 \left[ \begin{array}{c}
\frac{5850918398207460476949501506640014136017792143679}{25905291548014585554238994486937124864000000} \\
\frac{10317600297587322672417525639}{440141190696332728422480000} - \frac{22176170947759}{33508344870} \end{array} \right] \]
\[+ \frac{d^{\text{abcd}}_{A} d^{\text{abcd}}_{A}}{N_{R}} \left[ \begin{array}{c}
\frac{137762838063490156963463567256914633249848729}{68309157221399309555406722725916224000000} \\
\frac{15923437134502520764317446141}{2347419717137745515865600000} - \frac{59290512768143}{2843132292000} - \frac{40278295293983}{42646984380} \end{array} \right] \]
\[+ \frac{C_{2} A_{F}}{N_{R}} \left[ \begin{array}{c}
\frac{25891425129846636542816679113637834371}{920356792163320675817892446208000000} \\
\frac{22014640304042236864676686923}{16767283693841039399040000} - \frac{59290512768143}{189542152800} - \frac{1804393628665651}{625489104240} \end{array} \right] \]
\[+ \frac{C_{3} A_{F}}{N_{R}} \left[ \begin{array}{c}
\frac{1418362754727960165784082208324742517}{1284300051766864476020965376000000} \\
\frac{2649527312649305104376683}{3862475881756882790400} - \frac{59290512768143}{568626458400} - \frac{101176031536771}{63970476570} \end{array} \right] \]
\[+ n_{f} C_{2}^3 \left[ \begin{array}{c}
\frac{9440301206392517059551180859037562017787057}{31714965857292536579295784081711104000000} \\
\frac{88662795253}{193243050} + \frac{131795548}{218295} \end{array} \right] \]
\[+ n_{f}^2 C_{2}^2 \left[ \begin{array}{c}
\frac{776154252048358965658599007}{1424276065938472971760000} - \frac{1022665028447}{6087156075} \end{array} \right] \]
\[+ n_{f}^2 C_{A} C_{F} \left[ \begin{array}{c}
\frac{10327818039687232051823761}{398398899563313929280000} + \frac{1022665028447}{6087156075} \end{array} \right] \]
\[+ n_{f}^3 C_{F} \left[ \begin{array}{c}
\frac{391190356168864214011}{28457064254524235200} + \frac{1440968}{243243} \end{array} \right] \]
\[+ \frac{d_{A}^{abcd} d_{A}^{abcd}}{N_{R}} \left[ \begin{array}{c}
\frac{1062557701742420655559903}{47831059368392000000} + \frac{5928382591268105459}{8530540523505000} \end{array} \right] \]
\[+ \frac{C_{A} C_{F}}{N_{R}} \left[ \begin{array}{c}
\frac{4399336453016}{405810405} \end{array} \right] \]
\[+ n_{f} d_{A}^{abcd} d_{A}^{abcd} \left[ \begin{array}{c}
\frac{1444132331526475687660141}{276389512961562000000} + \frac{4789034376385918}{7456766191875} - \frac{1423427008}{945945} \end{array} \right] \]
\[, \]
\[\begin{align*}
&\frac{25752552164686212840197}{8621597950351848000} \zeta_3 + \frac{59290512768143}{1563722760600} \zeta_4 - \frac{128839202}{153153} \zeta_5 \\
&+ n_f C_A C_F \left[ -\frac{2893023190428781309670837247560962327027}{7247809738286150322065903013888000000} \zeta_3 - \frac{541630986863623}{1042481840400} \zeta_4 + \frac{64419601}{459459} \zeta_5 \right] \\
&+ \frac{63010566480259446207079}{57477319669001232000} \zeta_3 - \frac{1042481840400}{459459} \zeta_4 + \frac{1459459}{459459} \zeta_5 \\
&+ n_f^2 C_A^2 C_F \left[ -\frac{67436712195558046827798227902824541}{993803611447435941596860416000000} - \frac{1108384059230115101593813}{7663642622533497600} \zeta_3 \\
&+ \frac{1506311935054583}{3127445521200} \zeta_4 + \frac{4569981743}{7209972} \zeta_5 \right] \\
&+ n_f^2 C_F^2 \left[ \frac{149976730394272147964957809345185131}{258850347795933940073782250496000000} + \frac{249199939100703}{14073504845400} \zeta_3 \\
&+ \frac{64419601}{765765} \zeta_4 \right] \\
&+ \frac{n_f^2 C_A C_F}{4} \left[ \frac{9521719981900848384279388253327}{35492986123122712199887872000000} + \frac{249199939100703}{14073504845400} \zeta_3 - \frac{64419601}{765765} \zeta_4 \right] \\
&+ n_f^3 C_F \left[ \frac{543969361955458207942967542719}{380281994176314773570227200000} + \frac{128839202}{20675655} \zeta_3 \right] \\
&+ \frac{d_{abcd}^2 d_{abcd}^2}{N_R} \left[ \frac{13544672049501491403276515417}{530143748764348928000000} + \frac{379195537871486958507}{483002686294128000} \zeta_3 \\
&- \frac{5744913623184647}{469116828180} \zeta_5 \right] \\
&+ \frac{n_f^2 d_{abcd}^2 d_{abcd}^2}{N_R} \left[ \frac{6304830934247319711503473}{10843849400654350000000} + \frac{1361755723179569}{2020810952160} \zeta_3 \\
&- \frac{12596460722}{7810803} \zeta_5 \right]. \\
&\text{(B.8)}
\end{align*}\]

The corresponding results for the odd moments \(N = 3\) to \(N = 15\) read

\[\gamma_{\text{fit}}^{(3)}(N = 3) = \]

\[C_F^4 \left[ \frac{3416111945}{2239488} + \frac{24380}{81} \zeta_3 - \frac{2000}{3} \zeta_5 \right] \]

\[+ C_F^2 C_A \left[ \frac{40709323}{279936} - \frac{140057}{648} \zeta_3 - \frac{605}{3} \zeta_4 + \frac{2900}{3} \zeta_5 \right] \]

\[+ C_A^2 C_F^2 \left[ \frac{503877829}{559872} - \frac{4843}{27} \zeta_3 + \frac{605}{2} \zeta_4 - \frac{1325}{9} \zeta_5 \right] \]
\[+C_F C_A^2 \left[ \frac{72667541 + 125219}{69984} \cdot \frac{648}{6} \cdot \frac{605}{5950}{27} \cdot \zeta_3 - \frac{6}{3} \cdot \zeta_4 - \frac{27}{\zeta_5} \right] \]
\[+C_F^2 n_f \left[ \frac{38386673 + 3493}{279936} \cdot \frac{81}{3} \cdot \zeta_3 - \frac{110}{3} \cdot \zeta_4 - \frac{1000}{3} \cdot \zeta_5 \right] \]
\[+C_F^2 C_A n_f \left[ - \frac{22941613}{139968} + \frac{31547}{54} \cdot \zeta_3 - \frac{715}{3} \cdot \zeta_4 + \frac{500}{9} \cdot \zeta_5 \right] \]
\[+C_F C_A^2 n_f \left[ - \frac{2366971}{7776} - \frac{11483}{18} \cdot \zeta_3 + \frac{605}{3} \cdot \zeta_4 + \frac{7000}{27} \cdot \zeta_5 \right] \]
\[+C_F^2 n_f^2 \left[ \frac{1313443}{69984} - \frac{610}{9} \cdot \zeta_3 + \frac{100}{3} \cdot \zeta_4 \right] + C_F C_A n_f^2 \left[ \frac{79747}{5832} + \frac{610}{9} \cdot \zeta_3 - \frac{100}{3} \cdot \zeta_4 \right] \]
\[+C_F n_f^3 \left[ \frac{35587}{34992} + \frac{200}{81} \cdot \zeta_3 \right] + \frac{d_F^{abcd} d_F^{abcd}}{N_R} \left[ - \frac{85}{3} + \frac{7600}{9} \cdot \zeta_3 - \frac{7300}{9} \cdot \zeta_5 \right] \]
\[+n_f \frac{d_F^{abcd} d_F^{abcd}}{N_R} \left[ \frac{275}{3} + \frac{1960}{9} \cdot \zeta_3 - \frac{400}{9} \cdot \zeta_5 \right], \quad (B.9)\]

\[\gamma_3^{(3)}(N = 5) = \]
\[C_F^4 \left[ \frac{421526640437}{2187000000} + \frac{16940714}{25625} \cdot \zeta_3 - \frac{11032}{15} \cdot \zeta_5 \right] \]
\[+C_F^2 C_A \left[ - \frac{27591809}{1215000} - \frac{2335808}{10125} \cdot \zeta_3 - \frac{15554}{75} \cdot \zeta_4 + \frac{16408}{15} \cdot \zeta_5 \right] \]
\[+C_F^2 C_A^2 \left[ - \frac{151355831129}{218700000} + \frac{413533}{3750} \cdot \zeta_3 + \frac{7777}{25} \cdot \zeta_4 - \frac{6398}{9} \cdot \zeta_5 \right] \]
\[+C_F C_A^3 \left[ \frac{44149637147}{388800000} + \frac{4424851}{81000} \cdot \zeta_3 - \frac{7777}{75} \cdot \zeta_4 + \frac{13594}{135} \cdot \zeta_5 \right] \]
\[+C_F n_F C_A \left[ \frac{20671111851}{1093500000} + \frac{1181666}{10125} \cdot \zeta_3 - \frac{2828}{75} \cdot \zeta_4 - \frac{1456}{3} \cdot \zeta_5 \right] \]
\[+n_f C_F^2 C_A \left[ - \frac{8570931889}{36450000} + \frac{92291}{125} \cdot \zeta_3 - \frac{24262}{75} \cdot \zeta_4 + \frac{728}{9} \cdot \zeta_5 \right] \]
\[+n_f C_F C_A^2 \left[ - \frac{36596231437}{87480000} - \frac{2976818}{3375} \cdot \zeta_3 + \frac{21434}{75} \cdot \zeta_4 + \frac{50456}{135} \cdot \zeta_5 \right] \]
\[+n_f^2 C_F^2 \left[ \frac{269584781}{9112500} - \frac{67532}{675} \cdot \zeta_3 - \frac{278}{15} \cdot \zeta_4 \right] + n_f^2 C_F C_A \left[ \frac{44950469}{2430000} + \frac{67532}{675} \cdot \zeta_3 - \frac{728}{15} \cdot \zeta_4 \right] \]
\[+n_f C_F^3 \left[ - \frac{10064827}{109350000} - \frac{1146}{405} \cdot \zeta_3 + \frac{d_F^{abcd} d_F^{abcd}}{N_R} \right] \left[ \frac{38339}{80} + \frac{52857}{25} \cdot \zeta_3 - \frac{135968}{45} \cdot \zeta_5 \right] \]
\[+n_f \frac{d_F^{abcd} d_F^{abcd}}{N_R} \left[ \frac{74501}{450} + \frac{82432}{225} \cdot \zeta_3 - \frac{33152}{45} \cdot \zeta_5 \right], \quad (B.10)\]
\[ \gamma^{(3)}_{\text{ns}}(N = 7) = \]
\[ C_F \left[ \frac{8926031508626821967}{4098477035520000000} + \frac{107378934083}{3241350000} \zeta_3 - \frac{555362}{735} \zeta_5 \right] \]
\[ + C_A C_F^3 \left[ -2731357455838412101 + \frac{976102616819}{51861600000} \zeta_3 - \frac{1020151}{4900} \zeta_4 + \frac{832501}{735} \zeta_5 \right] \]
\[ + C_A^2 C_F^2 \left[ -1842858791204823727 + \frac{964491706751}{25930800000} \zeta_3 + \frac{3600453}{9800} \zeta_4 - \frac{3607811}{2940} \zeta_5 \right] \]
\[ + C_A^3 C_F \left[ 6856528500444857 \left( \frac{5974456320000}{44452800000} \zeta_3 - \frac{555362}{735} \zeta_5 \right) +\right] \]
\[ + n_f C_F^3 \left[ \frac{1462431578723251501}{65868380928000000} + \frac{1020151}{9800} \zeta_4 + \frac{942569}{2205} \zeta_5 \right] \]
\[ + n_f C_F^2 C_A \left[ -\frac{132713980134736771}{470488435200000} + \frac{447543373}{5292000000} \zeta_3 - \frac{1859803}{4900} \zeta_4 + \frac{2054}{21} \zeta_5 \right] \]
\[ + n_f C_A^2 C_F \left[ -661753424619487 + \frac{57952538719}{1344256720000} \zeta_3 + \frac{1674321}{4900} \zeta_4 + \frac{197849}{441} \zeta_5 \right] \]
\[ + n_f^2 C_F \left[ \frac{1463487948290143}{39207369600000} + \frac{8059127}{66150} \zeta_3 + \frac{2054}{35} \zeta_4 \right] \]
\[ + n_f^2 C_F C_A \left[ \frac{23660663137019}{1120210560000} + \frac{8059127}{66150} \zeta_3 - \frac{2054}{35} \zeta_4 \right] \]
\[ + n_f^3 C_F \left[ -\frac{902896393223}{840157920000} + \frac{4108}{945} \zeta_3 \right] \]
\[ + \frac{d_{\text{abcd}} d_{\text{abcd}}}{N_R} \left[ \frac{40269598361}{423360000} + \frac{885141073}{26460000000} \zeta_3 - \frac{3727489}{735} \zeta_5 \right] \]
\[ + \frac{n_f d_{\text{abcd}} d_{\text{abcd}}}{N_R} \left[ \frac{584326699}{226800000} + \frac{30610507}{66150000} \zeta_3 - \frac{142568}{147} \zeta_5 \right], \tag{B.11} \]

\[ \gamma^{(3)}_{\text{ns}}(N = 9) = \]
\[ C_F \left[ \frac{11729074138735897476581}{5041395218133760000000} + \frac{12756265517567}{393824025000} \zeta_3 - \frac{2166406}{2835} \zeta_5 \right] \]
\[ + C_A C_F^3 \left[ -\frac{1414388556845197343069}{8574651731520000000} - \frac{2630697602423}{1750329000000} \zeta_3 - \frac{13785409}{66150} \zeta_4 + \frac{4562074}{3969} \zeta_5 \right] \]
\[ + C_A^2 C_F^2 \left[ \frac{624559062347065092853}{1905478162560000000} + \frac{32816667064709}{5250987000000} \zeta_3 + \frac{13785409}{44100} \zeta_4 - \frac{91553}{54} \zeta_5 \right] \]
\[ + C_A^3 C_F \left[ \frac{44405787094076715779}{388873094000000000} - \frac{7312706287799}{300056400000} \zeta_3 - \frac{13785409}{132300} \zeta_4 + \frac{17432173}{23814} \zeta_5 \right] \]
\begin{align}
& + n_f C_F^2 \left[ \frac{29494412200734623467}{120045124241280000} + \frac{77319514799}{375070500} \zeta_3 + \frac{1253219}{33075} \zeta_4 - \frac{4180}{63} \zeta_5 \right] \\
& + n_f C_A C_F^2 \left[ - \frac{15104254419980130497}{47636954046000000} - \frac{230513485753}{250047000} \zeta_3 - \frac{9299719}{22050} \zeta_4 + \frac{20900}{189} \zeta_5 \right] \\
& + n_f C_A^2 C_F \left[ - \frac{2486710179097323617}{4536852768000000} - \frac{1091690689753}{937676250} \zeta_3 + \frac{25392719}{66150} \zeta_4 + \frac{4292332}{8505} \zeta_5 \right] \\
& + n_f^2 C_F^2 \left[ \frac{103302930942446363}{238184770320000} + \frac{41192947}{297675} \zeta_3 + \frac{9299719}{22050} \zeta_4 - \frac{4180}{63} \zeta_5 \right] \\
& + n_f^2 C_A C_F \left[ \frac{519982098287853}{22684263840000} + \frac{41192947}{297675} \zeta_3 - \frac{4180}{63} \zeta_4 \right] \\
& + n_f^3 C_F \left[ - \frac{40350728956471}{3402695760000} + \frac{8360}{1701} \zeta_3 \right] \\
& + \frac{d_{abcd} d_{abcd}}{N_R} \left[ \frac{2967911583758917}{416040608000000} + \frac{5621430297923}{12502350000} \zeta_3 - \frac{137680664}{19845} \zeta_5 \right] \\
& + \frac{d_{abcd} d_{abcd}}{N_R} \left[ \frac{29004062646461}{2700507600000} + \frac{11854062254}{22325625} \zeta_3 - \frac{3281344}{2835} \zeta_5 \right] \\
& + \gamma_{\text{ms}}^{(3)} (N = 11) = \\
& C_F^4 \left[ \frac{5942479928802007050870338007469}{245630685699637249824000000} + \frac{182985306975827551}{57659775502500} \zeta_3 - \frac{263657654}{343035} \zeta_5 \right] \\
& + C_A C_F^3 \left[ - \frac{1406439552456171828769629}{69047811800651376000000} - \frac{11683631163531071}{96099625833750} \zeta_3 - \frac{151689577}{727650} \zeta_4 \right] \\
& + \frac{2781069694}{2401245} \zeta_5 \\
& + C_A^2 C_F^2 \left[ - \frac{65080732936319895430465693}{33477726933649152000000} + \frac{1334072255443149799}{1537594013340000} \zeta_3 + \frac{151689577}{485100} \zeta_4 \right] \\
& - \frac{23116873}{10890} \zeta_5 \\
& + C_A^2 C_F \left[ \frac{1934103459065790006173}{1711041615360000000} - \frac{937411447527013}{24204549600000} \zeta_3 - \frac{151689577}{1455300} \zeta_4 \right] \\
& + \frac{14593135349}{14407470} \zeta_5 \\
& + n_f C_F^2 \left[ \frac{2807323127469698463268574221}{1063336301730031190400000} + \frac{118758665137891}{499218835500} \zeta_3 + \frac{151689577}{4002075} \zeta_4 \right] \\
& - \frac{502528}{693} \zeta_5 \right].
\end{align}
\begin{align}
+n_{f} C_{A} C_{F}^{2} & \left[ \frac{-28886000005383327025035503}{83694317343412288000000} + \frac{109012601223061}{110937519000} \zeta_{3} - \frac{1215792617}{2668050} \zeta_{4} \right] \\
+ \frac{251264}{2079} \zeta_{5} \\
+n_{f} C_{A}^{2} C_{F} & \left[ \frac{-1791290424092116907479}{3019275517104000000} - \frac{52497666373939}{41601569625} \zeta_{3} + \frac{3343998697}{8004150} \zeta_{4} \right] \\
+ \frac{37705252}{68607} \zeta_{5} \\
+n_{f}^{2} C_{F}^{2} & \left[ \frac{37084682543660792132933}{7671979088961264000000} - \frac{1093228621}{7203735} \zeta_{3} + \frac{251264}{3465} \zeta_{4} \right] \\
+n_{f}^{2} C_{A} C_{F} & \left[ \frac{31274848451808887}{1286623089675000} + \frac{1093228621}{7203735} \zeta_{3} - \frac{251264}{3465} \zeta_{4} \right] \\
+n_{f}^{3} C_{F} & \left[ \frac{633953354507911423}{498180460322160000} + \frac{502528}{93555} \zeta_{3} \right] \\
+ \frac{d_{F}^{abcd} d_{A}^{bcd}}{N_{R}} \left[ \frac{27544072006307986519}{15684548140800000} + \frac{112240841713507}{20170458000} \zeta_{3} - \frac{20755381868}{2401245} \zeta_{5} \right] \\
+n_{f} \frac{d_{F}^{abcd} d_{A}^{bcd}}{N_{R}} & \left[ \frac{550922957081846869}{1307045678400000} + \frac{104968283381}{180093375} \zeta_{3} - \frac{30026464}{22869} \zeta_{5} \right].
\end{align}

(\text{B.13})

\begin{align}
\gamma_{\text{ns}}^{(3)} (N = 13) =
+ C_{2}^{F} \left[ \frac{109254370623053143691943402173119101653}{440370321638438344304874602880000000} + \frac{10279814538281025664097}{32936416961252805000} \zeta_{3} \right. \\
- \frac{312960229682}{405810405} \zeta_{5} \\
+ C_{A} C_{2}^{F} \left[ \frac{-1374331496721755254616627043272450693}{598794143642010180678733920000000} - \frac{122111455215270203991}{1219867294861215000} \zeta_{3} \right. \\
- \frac{3662719609}{17567550} \zeta_{4} + \frac{42916930526}{38891855} \zeta_{5} \\
+ C_{A}^{2} C_{2}^{F} \left[ \frac{-448216384754306083916951583828253}{5583162178480281409492992000000} + \frac{6030472603025623852733}{548940286875467500} \zeta_{3} \right. \\
+ \frac{3662719609}{11711700} \zeta_{4} - \frac{680613539689}{270540270} \zeta_{5} \\
+ C_{A}^{3} C_{2}^{F} \left[ \frac{43589208948057316412721986370967}{39043092157204765069440000000} - \frac{29267653901342586703}{55836265244760000} \zeta_{3} \right. \\
- \frac{3662719609}{35135100} \zeta_{4} + \frac{281672586271}{221351130} \zeta_{5}. 
\end{align}
\[ + n_f C_F^3 \left[ \begin{array}{c}
\frac{7156992636760583450168754611387207}{256626061560615006517168000000} + \frac{2904432954831602777}{10967837815935055000000} \\
\frac{3662719609}{96621525} - \frac{6997864}{909} - \frac{3498932}{9509} \end{array} \right] \zeta_3
\]

\[ + n_f C_A C_F^2 \left[ - \frac{22054474323190222766265253462577}{598195947694315864813920000000} + \frac{7554888787686866551}{731189187729000000} \right] \zeta_3
\]

\[ - \frac{31181191789}{64414350} + \frac{3498932}{27027} \zeta_5 \]

\[ + n_f C_A^2 C_F \left[ - \frac{44016652745476286989105451999}{697198074235799376240000000} - \frac{10316777282946572197}{7677486471154500000} \right] \zeta_3
\]

\[ + \frac{86220020149}{1932403050} \zeta_4 + \frac{4331692556}{7378371} \zeta_5 \]

\[ + n_f^2 C_F^3 \left[ \frac{3742938630150542702886432311}{712138032969423648588000000} - \frac{19843204449}{1217431215} \zeta_3 + \frac{3498932}{540545} \zeta_4 \right]
\]

\[ + n_f^2 C_A C_F \left[ \frac{5068093117700672893019381}{19919944978165696464000000} - \frac{19843204449}{1217431215} \zeta_3 - \frac{3498932}{540545} \zeta_4 \right]
\]

\[ + n_f^3 C_F \left[ - \frac{19115924965400169467303}{14228532127261211760000} + \frac{6997864}{1216215} \right]
\]

\[ + \frac{d_{abcd}^p d_{abcd}^p}{N_R} \left[ \frac{107463452886020580200393}{510257562390576000000} + \frac{2581697503167399049}{393717254931000} \zeta_3 \right]
\]

\[ - \frac{4144734743192}{405810405} \zeta_5 \]

\[ + n_f d_{F}^{p b c d} d_{F}^{p b c d} \left[ \frac{625911243726496309379}{127564390597644000000} + \frac{27932248804801058}{4474059715125} \zeta_3 - \frac{3554796992}{2459457} \zeta_5 \right] \zeta_5 \]

\[ \gamma_{ns}^{(3)}(N = 15) = \]

\[ C_F \left[ \frac{1593418359971838203424224724210060722526639}{631314893100865210395468230688768000000} + \frac{259680235193827374263}{843172724208071808} \right] \zeta_3
\]

\[ - \frac{89608395343}{115945830} \zeta_5 \]

\[ + C_A C_F^3 \left[ - \frac{91674811575654880535423035274865836261}{3678991218536510500901412044800000000} - \frac{47105526692823873500839}{5621148494720478720000000} \right] \zeta_3
\]

\[ - \frac{205125530543}{983782800} \zeta_4 + \frac{34430154145}{29513484} \zeta_5 \]
\[ + C_A^2 C_F \left[ \frac{24794744696072980162780614700810631}{1225105030087795889113890816000000} + \frac{264518190199865408922107}{20075530382874240000} \zeta_3 \right. \\
\left. + \frac{205125530543}{655855200} \zeta_4 - \frac{415521270781}{144288144} \zeta_5 \right] \\
+ C_A^2 C_F \left[ \frac{11008793137870484173113743839424629}{99503159224449857776640000000000} - \frac{4285776141034052935649}{6551455122051840000} \zeta_3 \right. \\
\left. - \frac{205125530543}{1967565600} \zeta_4 + \frac{6701178363}{44270226} \zeta_5 \right] \\
+ n_f C_F^2 \left[ \frac{765704555141235451117809578252282141}{262785087038322812492958003200000} + \frac{5058205287659325397}{17548540505496000} \zeta_3 \right. \\
\left. + \frac{205125530543}{5410805400} \zeta_4 - \frac{7397890}{9009} \zeta_5 \right] \\
+ n_f C_A C_F^2 \left[ \frac{5957585274731644785588178652001289}{1531381626097448613923635200000000} + \frac{4197787411208679709}{3899675667888000} \zeta_3 \right. \\
\left. - \frac{1834288866343}{3607203600} \zeta_4 + \frac{3698945}{27027} \zeta_5 \right] \\
+ n_f^2 C_F^2 \left[ \frac{71853884960899868737940767607}{1081713375784028123313600000000} - \frac{16545594079801785051}{11699027000366400000} \zeta_3 \right. \\
\left. + \frac{5092615537943}{10821610800} \zeta_4 + \frac{77250091}{124740} \zeta_5 \right] \\
+ n_f^2 C_A C_F \left[ \frac{6735134366531635893107289}{25497529572052091473920000} + \frac{8409151022911}{48697248600} \zeta_3 - \frac{739789}{9009} \zeta_4 \right. \\
\left. + \frac{1024283569869843897775}{728500844915774042112} + \frac{1479578}{243243} \zeta_3 \right] \\
+ n_f^3 C_F \left[ - \frac{1024283569869843897775}{728500844915774042112} + \frac{1479578}{243243} \zeta_3 \right. \\
\left. - \frac{1022483569869843897775}{728500844915774042112} + \frac{1479578}{243243} \zeta_3 \right] \\
+ \frac{d^{abcd} d^{abcd}}{N_R} A \left[ \frac{328171179053663132153783483}{1350790928801233920000000} + \frac{145906081836810245387}{194983783394400000} \zeta_3 \right. \\
\left. - \frac{18940495337639}{1623241620} \zeta_5 \right] \\
+ n_f \frac{d^{abcd} d^{abcd}}{N_R} \left[ \frac{3421868696243684318010149}{61911250903389888000000} + \frac{22453298164820123}{34088074020000} \zeta_3 \right. \\
\left. - \frac{210864634}{135135} \zeta_5 \right], \quad (B.15)
\[ \gamma_{\text{ns}}^{(3)}(N = 3) = n_F C_F \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{72365 + 3200}{216} \zeta_3 \right] + n_J C_A \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{60629 - 1010}{648} \zeta_3 + \frac{200}{9} \zeta_5 \right] + n_J^2 \frac{d^{abc} d_{abc} N_R}{N} \left[ -\frac{1879}{324} \right], \tag{B.16} \]

\[ \gamma_{\text{ns}}^{(3)}(N = 5) = n_F C_F \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{1110626839 + 133952}{5467500} \zeta_3 \right] + n_J C_A \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{11071396 - 48608}{273375} \zeta_3 + \frac{1792}{45} \zeta_5 \right] + n_J^2 \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{150143}{273375} \right], \tag{B.17} \]

\[ \gamma_{\text{ns}}^{(3)}(N = 7) = n_F C_F \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{779352339134399 + 116051}{560105280000} \zeta_3 \right] + n_J C_A \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{4572181575853 - 3093067}{24004512000} \zeta_3 + \frac{2250}{49} \zeta_5 \right] + n_J^2 \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{13376649869}{12002256000} \right], \tag{B.18} \]

\[ \gamma_{\text{ns}}^{(3)}(N = 9) = n_F C_F \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{102127669986574697 + 85388204}{99243654300000} \zeta_3 \right] + n_J C_A \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{2755119043881017 - 1155263626}{31505922000000} \zeta_3 + \frac{19712}{405} \zeta_5 \right] + n_J^2 \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{767593737326}{738420046875} \right], \tag{B.19} \]

\[ \gamma_{\text{ns}}^{(3)}(N = 11) = n_F C_F \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{11160377330166434882669 + 39731371264}{13949052889024800000} \zeta_3 \right] + n_J C_A \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{4824168874878027761 - 103259298727}{164687755478400000} \zeta_3 + \frac{18200}{363} \zeta_5 \right] + n_J^2 \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{40350238520199037}{45289132756560000} \right], \tag{B.20} \]

\[ \gamma_{\text{ns}}^{(3)}(N = 13) = n_F C_F \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{18532792410668780136669601961 + 87687308676848}{28759420562226724269900000} \zeta_3 \right] + n_J C_A \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{4824168874878027761 - 103259298727}{164687755478400000} \zeta_3 + \frac{18200}{363} \zeta_5 \right] + n_J^2 \frac{d^{abc} d_{abc} N_R}{N} \left[ \frac{40350238520199037}{45289132756560000} \right]. \]
\[+n_j C_F \frac{d^{abc}d_{abc}}{N_R} \left[ \frac{33917531435390384588111467}{469267934581780417000000} - \frac{202143595520348}{4474059715125} \right] \]
\[+n_j C_A \frac{d^{abc}d_{abc}}{N_R} \left[ \frac{8444176392147052173457}{11173046061471143850000} \right], \quad (B.21)
\]
\[\gamma_{\text{ns}}^{(3)s} (N = 15) =
\[n_j C_F \frac{d^{abc}d_{abc}}{N_R} \left[ \frac{-12364402852607089809412352833493}{232027519105674032412672000000} + \frac{8992813421837}{185934942000} \right] \]
\[+n_j C_A \frac{d^{abc}d_{abc}}{N_R} \left[ \frac{-647090296462513369880498347}{2028212579595052730880000000} - \frac{17737091969947579}{4090568882400000} \right] \]
\[+n_j \frac{d^{abc}d_{abc}}{N_R} \left[ \frac{143754424718401819943003}{2228805032522035968000000} \right]. \quad (B.22)
\]

C. Time-like splitting function

Here we present the difference between the space- and time-like non-singlet splitting functions at four loops, defined by \(\delta P^{(3)\pm}(x) = P^{(3)\pm}_{\sigma=1}(x) - P^{(3)\pm}_{\sigma=-1}(x)\). The expression for \(\delta P^{(3)\pm}(x)\) reads

\[
\delta P^{(3)\pm}(x) = 16C_F \left\{ (1+x) \left( -792H_{-2}\varsigma_3 - 364H_{-2}\varsigma_2 - 64H_{-3}\varsigma_3 + 12H_{-4,0} + 768H_{-2,-1}\varsigma_2 - 28H_{-2,0} \\
+304H_{-2,2} + 768H_{-1,-2}\varsigma_2 - 792H_{-1,0}\varsigma_3 - 364H_{-1,0}\varsigma_2 + 608H_{-1,3} + 1584H_{-1,4} - 32H_{0,0}\varsigma_3 \\
-324H_{2,2} + 120H_{2,4} - 648H_{3,1} + 112H_{3,3} + 40H_{4,2} - 32H_{5,1} + 192H_{-4,-1,0} + 64H_{-3,-2,0} \\
-192H_{-3,-1,0} - 192H_{-2,-2,0} - 120H_{-2,-1,0} + 672H_{-2,-1,2} + 384H_{-2,2,1} - 192H_{-1,-3,0} \\
-120H_{-1,-2,0} - 672H_{-1,-2,2} + 768H_{-1,-1,0}\varsigma_2 - 1344H_{-1,-1,3} - 56H_{-1,0,0} - 1280H_{-1,0,0}\varsigma_2 \\
+304H_{1,2,0} + 384H_{1,2,2} + 912H_{1,3,0} + 768H_{1,3,1} - 64H_{0,0,0}\varsigma_2 - 324H_{2,1,0} + 56H_{2,3,0} \\
+48H_{3,2,0} + 764_{4,0,0} + 40H_{4,1,0} + 192H_{-2,-1,0} - 624H_{-2,-1,0} + 192H_{-1,-2,-1,0} \\
-624H_{-1,-2,0,0} + 192H_{-1,-1,-2,0} - 240H_{-1,-1,0,0} - 672H_{-1,-1,2,0} + 636H_{-1,0,0} \\
+384H_{-1,2,0,0} + 384H_{-1,2,1,0} + 24H_{2,2,0,0} + 116H_{3,0,0,0} + 48H_{3,1,0,0} + 384H_{-1,1,0,0} \\
-1296H_{-1,-1,0,0} + 704H_{-1,0,0,0} + 112H_{2,0,0,0} + 24H_{2,1,0,0} \right) (1-x) \left( 288H_{-3}\varsigma_3 + \frac{9}{2}H_0 \\
-393H_{2,-56}\varsigma_3 - 76H_{2,64}\varsigma_2 + 96H_{2,6}\varsigma_2 - 384H_{-3,-1}\varsigma_2 - 192H_{-2,-2}\varsigma_2 + 144H_{-2,0}\varsigma_3 \\
-288H_{-2,4} - 393H_{1,0} - 56H_{1,0}\varsigma_3 - 78H_{1,0}\varsigma_2 - 560H_{1,2} - 96H_{1,2}\varsigma_2 - 96H_{1,4} + 560H_{2,1} \\
-96H_{2,1}\varsigma_2 - 64H_{2,3} + 384H_{-3,-1,2} + 192H_{-2,-2,2} - 192H_{-2,-1,0}\varsigma_2 + 384H_{-2,-1,3} \\
+256H_{-2,0}\varsigma_2 - 96H_{-2,3,0} + 560H_{1,0,-64H_{1,0}\varsigma_2 - 560H_{1,1,0}\varsigma_2 - 96H_{1,1,0}\varsigma_2 - 64H_{1,3,0} \\
-48H_{2,2,0} + 96H_{-2,-2,0,0} + 192H_{-2,-1,2,0} + 16H_{1,0,0,0} + 288H_{-2,-1,0,0,0} - 64H_{-2,0,0,0,0} \\
-32H_{1,0,0,0,0} + 96H_{1,1,0,0},0) \right) \right\} - (20 + 264x)H_0\varsigma_3 - (28 - 268x)H_{-3,0} - (32 + 96x)H_{3,0}\varsigma_2 \]
\[-(32 + 224x)H_{-5,0} + \left(\frac{69}{2} + \frac{1641}{2}x\right)H_{0,0} - \left(\frac{77}{2} - \frac{725}{2}x\right)H_{0,0,0} - (46 + 306x)H_{0,0,1}\]
\[+ (48 - 1240x)H_{0,0,0,0} + (61 + 839x)H_{0,0,0,1} + (70 + 310x)H_{0,0,0,2} + (72 + 792x)H_{0,0,0,3}\]
\[+ (80 + 48x)H_{5,0} - (88 - 72x)H_{2,0,2} - (98 + 82x)H_{2,0,0,0} - (110 + 130x)H_{0,0,0,0}\]
\[-(112 - 80x)(H_{3,0,0,0,0} - (120 + 168x)H_{0,0,0,0,0} + (126 + 906x)H_{0,0,0,0,1}\]
\[-(144 + 112x)H_{0,0,0,0,2} + (168 - 56x)H_{0,\zeta,0,0,0} - (176 + 208x)(H_{3,2,0} + H_{3,1,0})\]
\[+ (184 + 480x)H_{-2,0,0} + (188 - 376x)H_{0,0} - \left(\frac{403}{2} + \frac{851}{2}x\right)H_{0,0,1} + (240 + 80x)H_6\]
\[-(244 + 276x)H_{3,0,0} + (248 + 1284x)H_{0,0,0} - (260 + 388x)H_{2,0,0} + (264 + 904x)H_5\]
\[-(280 + 88x)H_{-4,0,0} + (280 + 600x)H_{-3,0,0} - (296 - 56x)H_{-3,0,0,0} + (300 - 61x)H_{-3,-1,0}\]
\[-(320 - 20x)H_{-4,0,0} - (342 + 550x)H_{1,0,0} + (352 + 999x)H_{-3,2,0} + (368 + 688x)H_{-2,2,0}\]
\[-(378 + 1506x)H_{1,0,1} - (400 - 720x)H_{3,0,2} - (400 - 304x)H_{-3,0,0,2} - (416 - 352x)H_{-3,0,3}\]
\[+ (432 - 144x)H_{-4,2} - (432 + 720x)H_{4,1} + (442 + 834x)H_{3,0} - (448 + 1088x)H_{-3,2}\]
\[-(480 - 640x)H_{2,0,0} + (584 + 1064x)H_{-2,0,0,0} - (864 + 1184x)H_{-2,0,0,0} + (880 + 1520x)H_{-2,0}\]
\[+ p_{\eta,\zeta}(-x)\left(-456H_{-4,0} - 1248H_{-3,0} + 192H_{-3,2} + 910H_{-2,4} + 96H_{-2,3} + 36H_{-2,2}\right)\]
\[-24H_0 + \left(\frac{45}{2} - \frac{39}{2}H_{0,0} - 308H_{0,0,0} + 24H_3 - 48H_{3,0} + 72H_{3,2} + 54H_4\right)\]
\[-120H_4 - 2H_5 - 160H_6 - 162H_{-4,0} + 52H_{-4,0} + 1152H_{-3,0}\]
\[-124H_{-3,0,0} + 12H_{-3,0,0} + 1312H_{-3,0} + 1152H_{-2,0} - 124H_{-2,0} - 124H_{-2,0,0} - 192H_{-2,0}\]
\[-15H_{-2,0} - 96H_{-2,0,0} + 216H_{-2,0,0} - 36H_{-2,0} - 16H_{-2,0} - 240H_{-2,0} - 172H_{-2}\]
\[+ 1152H_{-1,0} + 1248H_{-1,0} - 3H_{-1,0} + 12H_{-1,0} + 910H_{-1,0} + 36H_{-1,0}\]
\[-72H_{1,0} - 32H_{1,0} + 144H_{-1,0} + 1088H_{-1,0} - 180H_{1,0} + 612H_{0,0,0} + 36H_{0,0,1} + 12H_2\]
\[-24H_2 + 36H_{2,0} + 96H_{2,0} + 18H_{3,0} + 136H_{3,0} + 192H_{3,0} + 48H_{3,0} + 12H_4\]
\[-144H_4 - 96H_5 - 224H_6 + 384H_6 + 144H_6 + 96H_6 + 40H_6 + 48H_6\]
\[-1152H_{-3,0} - 12H_{-3,0} - 72H_{-3,0} + 56H_{-3,0} + 26H_{-3,0} - 1152H_{-2,0}\]
\[-1152H_{-2,0} - 166H_{-2,0} + 166H_{-2,0} + 12H_{-2,0} + 212H_{-2,0} - 18H_{-2,0}\]
\[-1408H_{-2,0} + 8H_{-2,0} - 8H_{-2,0} + 8H_{-2,0} + 1056H_{-3,0} + 1152H_{-2,3,1}\]
\[+ 144H_{-1,0} - 1152H_{-2,0} + 1152H_{-2,0} + 1152H_{-2,0} + 1152H_{-2,0} + 1152H_{-2,0}\]
\[+ 2880H_{-1,0} + 30H_{-1,0} + 672H_{-1,0} + 240H_{-1,0} - 36H_{-1,0} - 16H_{-1,0}\]
\[+ 96H_{-1,0} + 48H_{-1,0} + 48H_{-1,0} + 48H_{-1,0} + 48H_{-1,0} + 48H_{-1,0} + 48H_{-1,0}\]
\[-72H_{0,0,0} + 14H_{0,0,0} + 48H_{3,0} + 48H_{3,0} + 48H_{3,0} + 48H_{3,0} + 48H_{3,0}\]
\[-48H_{3,0,0} + 48H_{3,0,0} + 576H_{2,0,0} + 1152H_{2,0,0} + 1152H_{2,0,0} + 1152H_{2,0,0}\]
\[-960H_{-2,0,0} - 768H_{-2,0,0} + 144H_{-2,0,0} + 352H_{-2,0,0} + 352H_{-2,0,0} + 352H_{-2,0,0}\]
\[+ 1152H_{-1,2,2} + 1152H_{-1,2,2} + 768H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2}\]
\[-1152H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2}\]
\[-1152H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2} + 1152H_{-1,2,2}\]
\[-1728 \zeta_{-1,-1,3,0} - 1536 \zeta_{-1,-1,3,0} - 54 \zeta_{-1,0,0,0} - 960 \zeta_{-1,0,0,0} \zeta_2 + 96 \zeta_{-1,2,0,0} + 96 \zeta_{-1,2,1,0} + 320 \zeta_{-1,3,0,0} + 384 \zeta_{-1,1,0,0} + 64 \zeta_{0,0,0,0} + 48 \zeta_{0,3,0,0} + 576 \zeta_{-2,-1,1,0} - 1120 \zeta_{-1,-1,0,0} + 576 \zeta_{-1,-1,-2} - 1120 \zeta_{-1,-1,-2,0} + 1152 \zeta_{-1,-1,-2,0} + 288 \zeta_{-1,-1,-2,0} - 768 \zeta_{-1,-1,-2,0} - 768 \zeta_{-1,-1,-2,0} - 96 \zeta_{-1,2,0,0} + 60 \zeta_{0,0,0,0} + 1728 \zeta_{-1,-1,-1,0,0} - 1024 \zeta_{-1,-1,-1,0,0} - 80 \zeta_{-1,0,0,0} + 192 \zeta_{0,0,0,0} \right) \]

\[
+ p_{u_0}(x) \left( -264 \zeta_{-4} \zeta_2 - 336 \zeta_{-3} \zeta_3 + 72 \zeta_{-3} \zeta_2 + \frac{29}{8} H_0 + 144 H_0 \zeta_5 - \frac{45}{2} H_0 \zeta_4 + 25 H_0 \zeta_3 + \frac{9}{2} H_0 \zeta_2 - 124 H_0 \zeta_2 \zeta_3 - 66 H_2 \zeta_3 + 24 H_2 \zeta_3 - 64 H_3 \zeta_3 + 15 H_4 + 24 H_4 \zeta_2 + 96 H_5 - 320 H_6 + 192 H_{-5,0} - 36 H_{-4,0} + 48 H_{-4,2} + 384 H_{-3,0} \zeta_2 - 280 H_{-3,3} + 128 H_{-3,3} + 192 H_{-2,2} - 2 \zeta_2 - 168 H_{-2,0} \zeta_3 + 36 \zeta_{-2} \zeta_2 + 144 \zeta_{-2,4} + 76 H_{0,0} \zeta_4 - 280 H_{0,0} \zeta_4 + 32 H_{1,-3} \zeta_2 - 66 H_{1,0} \zeta_4 + 24 H_{1,0} \zeta_3 + 144 H_{1,4} - 256 H_{1,5} + 16 H_{2,-2} \zeta_2 - 32 H_{2,0} \zeta_3 - 64 H_{2,1} \zeta_3 + 96 H_{2,3} - 256 H_{2,4} + 13 H_{3,0} + 72 H_{3,0} \zeta_2 + 24 H_{2,2} - 224 H_{2,3} + 60 H_{4,0} - 72 H_{4,1} - 224 H_{4,2} + 288 H_{5,0} - 320 H_{5,1} - 432 H_{4,-1} + 336 H_{4,-0} - 336 H_{3,-2} - 0 + 144 H_{3,-3,1} - 192 H_{3,-3,1} - 192 H_{3,-2,2} - 192 H_{3,-2,1} - 1 - 192 H_{-2,1,3} - 192 H_{-2,0,0} \zeta_2 + 48 H_{-2,3,3} + \frac{45}{4} H_0,0,0,0 + 160 H_{0,0,0} \zeta_4 + 36 H_{0,0,0} \zeta_2 + 144 H_{-1,-4,0} + 96 H_{-1,-3,0} - 128 H_{-1,-3,2} + 16 H_{-1,-2,0} \zeta_2 - 128 H_{-1,-2,3} - 64 H_{1,1,0} \zeta_3 - 192 H_{1,1,4} + 32 H_{2,0} \zeta_2 - 128 H_{1,2,3} + 96 H_{1,3,0} - 160 H_{1,3,2} - 304 H_{1,4,0} - 288 H_{1,4,1} + 48 H_{2,2} - 3,0 + 48 H_{2,2} - 24 H_{2,2} - 249 H_{2,2} + 112 H_{2,0} \zeta_2 + 32 H_{2,1} \zeta_2 - 128 H_{2,3,3} + 72 H_{2,2,0} - 96 H_{2,2,2} - 240 H_{2,3,3} - 160 H_{2,3,1} + 120 H_{3,0,0} + 24 H_{3,1,0} - 96 H_{3,1,2} - 192 H_{3,2,0} - 96 H_{3,2,1} - 240 H_{4,0,0} - 224 H_{4,1,0} - 96 H_{4,1,1} + 384 H_{4,-1,1} - 1 - 576 H_{-3,1,0} + 336 H_{3,0,0} + 192 H_{-2,1,0} - 336 H_{2,-2,0} + 144 H_{-2,2,0,0} - 192 H_{-2,1,0} - 96 H_{-2,1,2,0} - 108 H_{-2,2,0} - 70 H_{0,0,0,0} + 256 H_{0,0,0,0} \zeta_2 - 192 H_{-1,-3,1,0} + 128 H_{-1,-3,0,0} - 96 \zeta_{-1,2,0} - 96 H_{1,2,0,0} - 64 H_{1,2,2,0} - 93 H_{1,0,0,0} + 144 H_{1,0,0,0} \zeta_2 + 64 H_{1,1,0,0} \zeta_2 - 128 H_{1,1,3,0} + 96 H_{1,2,2,0} - 96 H_{2,2,0,0} - 96 H_{2,2,1,0} - 176 H_{3,0,0,0} - 96 H_{3,1,0,0} - 96 H_{3,1,1} + 384 H_{2,1,0,0} - 432 H_{2,0,0,0} + 192 H_{2,0,0,0} + 180 H_{0,0,0,0} - 192 H_{1,-2,1,0} + 48 H_{1,-2,0,0} + 144 H_{1,0,0,0,0} + 144 H_{1,0,0,0,0} - 16 H_{1,2,0,0,0} - 144 H_{1,0,0,0,0} - 16 H_{2,1,0,0,0} - 272 H_{0,0,0,0,0,0} - 64 H_{1,1,1,1,0,0,0} + 192 H_{1,1,1,1,0,0,0} \right) \]

\[
+ 16 C_3^2 C_A \left( \left( 1 + x \right) \frac{864 \zeta_{-2} \zeta_3 + \frac{874}{3} H_{-2} \zeta_3 - 32 H_3 \zeta_3 - 384 \zeta_{-2,1} \zeta_2 + 238 H_{-2} - 288 H_{-2,2} - 864 \zeta_{-1,2} \zeta_2 + 864 \zeta_{-1,2} \zeta_2 + \frac{874}{3} H_{-1,2} - 576 H_{-1,3} - 1704 H_{-1,4} - 16 H_{2,0} \zeta_3 + \frac{1016}{3} H_{2,2} - 16 H_{2,2} + 284 H_{2,4} + \frac{2032}{3} H_{3,1} - 32 H_{3,1} \zeta_2 - 96 H_{3,3} + 48 H_{4,2} + 192 H_{5,1} + \frac{20}{3} H_{2,1,1,0} + 752 H_{-2,2,1} + 224 H_{-1,3,0} + \frac{20}{3} H_{-1,2,0} + 752 H_{-1,2,2} - 864 H_{-1,1,0} \zeta_2 \right) \right)
\[
\begin{align*}
+ \left( \frac{1492}{3} + \frac{3692}{3} x \right) H_{-3} \zeta_{2} + \left( \frac{1037}{2} - \frac{1251}{2} x \right) H_{2,0} + (768+656x)H_{-2,-1,0,0} \\
-(782+1114x)H_{-2,0,0,0} - \left( \frac{2776}{3} + \frac{5144}{3} x \right) H_{-2,3} + \left( \frac{2810}{3} + \frac{3910}{3} x \right) H_{-2,0} \zeta_{2} \\
-48 \pi H_{0,0,0,0,0,0,0,0} + 140x H_{0,0,0,0,0,0,0,0} + p_{qq}(-x) \left( 660H_{-4} \zeta_{2} + 1456H_{-3} \zeta_{3} - 288H_{-3} \zeta_{2} \\
-835H_{-2} \zeta_{4} - \frac{136}{3} H_{-2} \zeta_{3} - \frac{1258}{3} H_{-2} \zeta_{2} + 36H_{0} \zeta_{5} + \frac{65}{12} H_{0} \zeta_{4} + \frac{2155}{6} H_{0} \zeta_{3} - \frac{79}{4} H_{0} \zeta_{2} \\
+390H_{0} \zeta_{2} \zeta_{3} - 36H_{3} + 72H_{3} \zeta_{3} - 60H_{3} \zeta_{2} - 361H_{4} + 156H_{4} \zeta_{2} - 168H_{5} + 320H_{5} + 96H_{5} \right) \\
+194H_{-4,0} - 696H_{-4,2} + 1344H_{-3,1} \zeta_{2} + \frac{940}{3} H_{-3,0} + 1468H_{-3,0} \zeta_{2} + \frac{688}{3} H_{-3,2} \\
-1584H_{-3,3} + 1344H_{-2,-2} \zeta_{2} - 1456H_{-2,-1} \zeta_{3} + \frac{512}{3} H_{-2,-1} \zeta_{2} - \frac{151}{2} H_{-2,0} + 1184H_{-2,0} \zeta_{3} \\
-\frac{956}{3} H_{-2,0} \zeta_{2} + \frac{722}{3} H_{-2,2} + 8H_{-2,2} \zeta_{2} + \frac{582}{3} H_{-2,3} + \frac{2024}{3} H_{-2,4} - 1344H_{-1,3} \zeta_{2} \\
-1456H_{-1,-2} \zeta_{3} + \frac{512}{3} H_{-1,-2} \zeta_{2} - 835H_{-1,0} \zeta_{4} + \frac{136}{3} H_{-1,0} \zeta_{3} - \frac{1258}{3} H_{-1,0} \zeta_{2} + 144H_{-1,3} \zeta_{3} \\
+16H_{-1,3} \zeta_{2} + 184H_{-1,4} - 1376H_{-1,5} + 27H_{0,0} + 454H_{0,0} \zeta_{4} + \frac{316}{3} H_{0,0} \zeta_{3} + \frac{1138}{3} H_{0,0} \zeta_{2} \\
-18H_{2,0} + 36H_{2,0} \zeta_{3} - 30H_{2,0} \zeta_{2} - 112H_{2,2} \zeta_{2} + \frac{361}{3} H_{3,0} + 164H_{3,0} \zeta_{2} + 224H_{3,1} \zeta_{2} + \frac{248}{3} H_{3,2} \\
+38H_{4,0} + 248H_{4,1} + 112H_{4,2} + 288H_{5,0} + 448H_{5,1} - 72H_{-4,1,0} - 264H_{-4,0,0} - 24H_{-3,2,0} \\
- \frac{352}{3} H_{-3,-1,0} + 1344H_{-3,-1,2} + 420H_{-3,0,0} - 680H_{-3,2,0} + 680H_{-3,2,0} - 896H_{-3,2,1} - 24H_{-2,3,0} \\
- \frac{352}{3} H_{-2,-2,0} + 1344H_{-2,-2,2} + 1344H_{-2,-1,1} \zeta_{2} + \frac{1072}{3} H_{-2,-1,0} - 1920H_{-2,-1,0} \zeta_{2} \\
- \frac{512}{3} H_{-2,-1,2} + 246H_{-2,-1,3} + \frac{1072}{3} H_{-2,0,0} + 1648H_{-2,0,0} \zeta_{2} + \frac{184}{3} H_{-2,2,0} - \frac{496}{3} H_{-2,2,1} \\
- 448H_{-2,2,2} - 1232H_{-2,3,0} - 1344H_{-2,3,1} - 72H_{-1,-4,0} - \frac{352}{3} H_{-1,-3,0} + 1344H_{-1,-3,2} \\
+ 1344H_{-1,-2,1} \zeta_{2} + \frac{1072}{3} H_{-1,-2,0} - 1920H_{-1,-2,0} \zeta_{2} - \frac{512}{3} H_{-1,-2,2} + 224H_{-1,-2,3} \\
+ 1344H_{-1,-1,2} \zeta_{2} - 1456H_{-1,-1,0} \zeta_{3} + \frac{512}{3} H_{-1,-1,0} \zeta_{2} - \frac{1024}{3} H_{-1,-1,3} + 134H_{-1,-1,4} \\
- 151H_{-1,0,0} + 134H_{-1,0,0} \zeta_{3} - \frac{1048}{3} H_{-1,0,0} \zeta_{2} + \frac{722}{3} H_{-1,2,0} + 8H_{-1,2,0} \zeta_{2} - \frac{496}{3} H_{-1,2,2} \\
- \frac{136}{3} H_{-1,3,0} + \frac{992}{3} H_{-1,3,1} - 448H_{-1,3,2} - 1016H_{-1,4,0} - 1344H_{-1,4,1} + \frac{345}{4} H_{0,0,0} \\
- 192H_{0,0,0} \zeta_{3} + 304H_{0,0,0} \zeta_{2} + 112H_{2,0,0} \zeta_{2} + 112H_{2,1,0} \zeta_{2} - \frac{160}{3} H_{3,0,0} + \frac{248}{3} H_{3,1,0} + 64H_{4,0,0} \\
+ 112H_{4,0,0} + 624H_{-3,-1,0,0} - 664H_{-3,0,0,0} + 672H_{-2,-2,0} - 1344H_{-2,-1,1,2} \\
- 320H_{-2,-1,0,0} + 1120H_{-2,-1,2,0} + 896H_{-2,-1,2,1} + 572H_{-2,0,0,0} - 416H_{-2,2,0} \\
- 448H_{-2,2,1,0} + 624H_{-1,-3,0,0} - 1344H_{-1,-2,1,2} - 320H_{-1,-2,0,0} + 1120H_{-1,-2,2,0} \\
+ 896H_{-1,-2,2,1} - 1344H_{-1,-1,2,2} + 1344H_{-1,-1,0} \zeta_{2} - 268H_{-1,-1,0,0} \zeta_{2} + \frac{136}{3} H_{-1,-1,0,0} \\
- 249H_{-1,-1,0,0} \zeta_{2} - \frac{512}{3} H_{-1,-1,2,0} + 896H_{-1,-1,2,2} + 2016H_{-1,-1,1,3,0} + 1792H_{-1,-1,3,1}
\end{align*}
\]
\[ +765H_{-1,0,0,0} + 1200H_{-1,0,0,0} \zeta_2 - \frac{320}{3} H_{-1,2,0,0} - \frac{496}{3} H_{-1,2,1,0} - 384H_{-1,3,0,0} - 48H_{-1,3,1,0} \]
\[-700 \frac{H_{0,0,0,0} - 224H_{0,0,0,0} \zeta_2 - 48H_{3,0,0,0} - 672H_{-2,1,0,0} + 1360H_{-2,1,0,0}}{3} \]
\[-744H_{-2,0,0,0} - 672H_{-1,2,1,0} - 1360H_{-2,0,0,0} - 672H_{-1,2,1,0} - 1,2,0 \]
\[-1344H_{-1,1,2,0} - 608H_{-1,1,0,0} + 896H_{-1,1,2,0} + 896H_{-1,1,2,0} + 544H_{-1,0,0,0} \]
\[+96H_{-1,2,0,0} - 310H_{0,0,0,0} - 2016H_{-1,1,1,0,0} + 1408H_{-1,1,0,0,0} - 200H_{1,0,0,0,0} \]
\[-48H_{0,0,0,0,0} \] + \( p \phi(x) \left( \frac{84H_{-1,0,0,0} + 280H_{-3,0,0,0} + 36H_{-3,0,0,0} \zeta_2 + \frac{587}{48} H_{0,0} - 26H_{0,0} \zeta_3 + \frac{149}{4} H_{0} \zeta_4}{4} \right) \]
\[-\frac{200}{3} H_{0} \zeta_3 - \frac{229}{4} H_{0} \zeta_2 - 6H_{6} \zeta_2 \zeta_3 + 313H_{2} \zeta_4 - 44H_{2} \zeta_3 + 320H_{3} \zeta_3 + 226H_{4} + 12H_{4} \zeta_2 \]
\[+ \frac{640}{3}H_{5} + 40H_{6} - 96H_{5,0} - 90H_{4,0} + 120H_{4,2} - 384H_{3,0,0} - 156H_{3,0,0} \zeta_2 \]
\[-96H_{3,2} - 24H_{3,3} - 192H_{2,2,0} - 2\zeta_2 + 140H_{2,0,0} \zeta_3 + 18H_{2,0,0} \zeta_3 - 96H_{2,3,0} - 24H_{2,4} \]
\[+ \frac{11}{4} H_{0,0} + 34H_{0,0,0,0} \zeta_4 - \frac{356}{3} H_{0,0,0,0} \zeta_3 - \frac{464}{3} H_{0,0,0,0} \zeta_2 - 272H_{1,1,0} \zeta_3 + 313H_{1,0} \zeta_4 + 44H_{1,0} \zeta_3 \]
\[+ \frac{320}{3} H_{1,2} \zeta_3 - 232H_{1,4} + 64H_{1,5} - 136H_{2,2} - 2\zeta_2 + 352H_{2,0,0} \zeta_3 + \frac{52}{3} H_{2,0} \zeta_2 + 320H_{2,1} \zeta_3 \]
\[-\frac{232}{3} H_{2,3} + 112H_{2,4} + \frac{226}{3} H_{3,0,0} + 4H_{3,0,0} \zeta_2 + 48H_{3,0,0} + 64H_{3,3} + \frac{508}{3} H_{4,0} + 144H_{1,1} \]
\[+ 32H_{4,2} + 72H_{5,0} + 128H_{5,1} + 408H_{4,1,0} - 192H_{4,0,0} + 360H_{3,1,0} - 120H_{3,0,0} \]
\[+ 160H_{3,2,0} + 24H_{3,2,0} + 140H_{3,3,0} + 128H_{2,3,0} + 224H_{2,3,0} - 60H_{2,0,0} \]
\[+ 80H_{2,2,0} - 192H_{2,1,0} \zeta_2 + 160H_{2,2,0} - 136H_{2,0,0} \zeta_2 - 48H_{2,2,0} \]
\[+ 64H_{2,2,2} + 56H_{2,3,0} + 128H_{2,3,1} - \frac{597}{4} H_{0,0,0,0} + 120H_{0,0,0,0} \zeta_3 - 250H_{0,0,0,0} \zeta_2 + 120H_{1,1,4,0} \]
\[-240H_{1,1,0} - 320H_{1,1,2,0} - 316H_{1,1,2,0} - 320H_{1,1,2,0} + 384H_{1,0,0} \zeta_3 + \frac{104}{3} H_{1,1,0} \zeta_2 \]
\[+ 320H_{1,1,0} \zeta_3 + 24H_{1,1,1} + 164H_{1,2,2} - \frac{232}{3} H_{1,3,0} + 128H_{1,3,2} + 208H_{1,4,0} + 384H_{1,4,1} \]
\[+ 168H_{2,1,0,0} - 120H_{2,1,0,2} + 160H_{2,2,0} - 2,0 + \frac{400}{3} H_{2,0,0,0} + 8H_{2,0,0} \zeta_2 + 64H_{2,1,3} + 96H_{2,3,0} \]
\[+ 128H_{2,3,1} + 128H_{2,2,0} + 48H_{3,0,0} + 48H_{3,1,0} - 16H_{4,0} + 32H_{4,1,0} - 448H_{3,1,0} + 32H_{3,1,0} - 448H_{3,1,0} \]
\[+ 576H_{3,1,0} - 224H_{3,0,0} - 224H_{2,2,2} - 2,1,0 + 376H_{2,2,2} - 2,0,0 - 224H_{2,2,2} - 2,1,0 \]
\[-120H_{2,2,2} - 2,1,0 + 80H_{2,2,2} - 2,1,0 + 18H_{2,2,2} - 0,0,0 + 64H_{2,2,2} - 0,0,0 + 64H_{2,2,2} - 0,0,0 + 748 H_{0,0,0,0} \]
\[-64H_{0,0,0,0} \zeta_2 + 96H_{1,3,0} + 192H_{1,3,0} - 3,0 + 48H_{1,1,0} - 2,0 - 240H_{1,2,0} - 2,0 + 160H_{1,1,2,0} \]
\[+ 400H_{1,0,0,0} + 24H_{1,0,0,0} \zeta_2 + 256H_{1,1,1} - 3,0 + 64H_{1,1,3} + 128H_{1,2,0} - 2,0 - \frac{176}{3} H_{1,2,0,0} \]
\[-64H_{1,3,0} + 128H_{1,3,1,0} + 48H_{2,2,2} - 2,0,0 + 184H_{2,2,0} - 2,0,0 - 8H_{2,0,0,0} + 128H_{2,2,0,0} \]
\[-176 \frac{H_{2,1,0,0} - 192H_{2,2,0,0} - 176H_{3,0,0} - 192H_{3,0,0} - 448H_{2,1,0,0} - 192H_{2,2,0,0} + 546H_{2,2,0,0} \]
\[-160H_{2,2,0,0} + 310H_{0,0,0,0} + 96H_{1,2,1,0} + 168H_{1,2,0,0} + 144H_{1,1,0,0} \]
\[+ 256H_{1,1,2,0} - 176H_{1,1,2,0} - 192H_{1,1,2,0} + 416H_{1,2,0,0} - 192H_{1,2,1,0} - 296H_{2,0,0,0} \]
\[-416H_{2,1,0,0} + 192H_{2,1,0,0} + 48H_{0,0,0,0} + 200H_{1,0,0,0} - 512H_{1,1,0,0} \]
\[ -576H_{1,1,1,0,0,0} \) \]

\[ + \frac{16}{3} C_f^2 C_A \left\{ (1+x) \left( -708H_{-2} \xi_3 - 164H_{-2} \xi_2 + 120H_3 \xi_3 + 720H_{-2,-1} \xi_2 - 336H_{-2,0} \right) \\
+ 204H_{-2,2} + 720H_{-1,-2} \xi_2 - 708H_{-1,0} \xi_3 - 164H_{-1,0} \xi_2 + 408H_{-1,2} + 136H_{-1,4} + 40H_2 \xi_3 \\
- 336H_{2,2} + 24H_{2,2} \xi_2 + 162H_{2,4} - 672H_{3,1} + 48H_{3,1} \xi_2 - 144H_{3,2} + 108H_{3,3} - 432H_{4,1} \\
- 36H_{4,2} - 144H_{5,1} + 80H_{-2,-1,0} - 624H_{-2,-2,0} + 288H_{-2,2,1} - 192H_{-1,-3,0} + 80H_{-1,-2,0} \\
- 624H_{-1,-2,2} + 720H_{-1,-1,0} \xi_2 - 124H_{-1,-1,0} - 72H_{-1,0,0} - 1104H_{-1,0,0} \xi_2 + 204H_{-1,2,0} \\
+ 288H_{-1,2,2} + 744H_{-1,3,0} + 576H_{-1,3,1} + 96H_{-3,-3,0} - 36H_{2,0,0} \xi_2 - 336H_{2,1,0} + 24H_{2,1,0} \xi_2 \\
+ 54H_{2,3,0} + 96H_{3,1,0} - 99H_{4,0,0} - 36H_{4,1,0} + 192H_{-2,-1,0} + 192H_{-1,-1,0} \\
- 600H_{1,-2,0} + 192H_{1,-1,-2,0} + 160H_{1,-1,0,0} - 624H_{1,-1,2,0} - 78H_{1,-1,0,0} \\
+ 288H_{1,-2,0,0} + 288H_{1,-2,1,0} + 96H_{2,-2,0,0} - 42H_{2,2,0,0} - 99H_{3,3,0,0} - 84H_{3,1,0,0} \\
+ 384H_{-1,-1,-1,0,0} - 1224H_{-1,-1,0,0,0} + 768H_{-1,0,0,0,0} - 72H_{2,0,0,0,0} - 126H_{2,1,0,0,0} \right) \\
+ (1-x) \left( 360H_{-4} \xi_2 + 372H_{-3} \xi_3 - \frac{16997}{27} H_0 - \frac{9800}{9} H_2 - 276H_3 \xi_3 + 40H_2 \xi_2 - 24H_{-5,0} \right) \\
- 480H_{-3,-1} \xi_2 + 408H_{-3,0} \xi_2 - 240H_{-2} \xi_2 + 186H_{-2,0} \xi_3 - 324H_{-2,4} - \frac{9800}{9} H_{1,0} \\
- 276H_{1,0} \xi_3 + 40H_{1,0} \xi_2 - 528H_{1,2} - 96H_{1,2} \xi_2 - 144H_{1,4} - 528H_{2,1} - 96H_{2,1} \xi_2 - 48H_{2,3} \\
+ 432H_{-3,-1,2} + 48H_{-3,-3,0} + 216H_{-2,-2,2} - 240H_{-2,-1,0} \xi_2 + 432H_{-2,-1,3} + 288H_{-2,0,0} \xi_2 \\
- 108H_{-2,3,0} - 192H_{1,-3,0} - \frac{2200}{3} H_{1,0,0,0} - 96H_{1,0,0} \xi_2 - 528H_{1,1,0} - 96H_{1,1,0} \xi_2 - 48H_{1,3,0} \\
- 96H_{2,0,0} - 96H_{3,-1,-1,0} + 48H_{2,-2,1,0} + 180H_{2,-2,0,0} + 48H_{2,-1,-2,0} \\
+ 216H_{-2,-2,0} - 192H_{1,-2,0,0} + 228H_{1,0,0,0} + 144H_{1,2,0,0} + 144H_{1,2,0,0,0} + 96H_{2,-1,0,0} \\
+ 396H_{2,-1,0,0,0} - 192H_{2,-2,0,0,0} + 384H_{1,0,0,0,0} + 432H_{1,1,0,0,0} \right) + (10-346x)H_{-3,0} \\
- (16+372x)H_{-2,0,0} - \left( \frac{87}{2} + \frac{21}{2} x \right) H_{0,0,4} - (45 + 105x) H_{5,0} + \left( \frac{123}{2} + \frac{1001}{2} x \right) H_{0,0,3} \\
+ (66+930x) H_{0,0,4} - (75+51x) H_{0,0,5} - (84-132x) H_{3,2,0,0} + (90-174x) H_{0,0,2} \xi_3 \\
+ \left( \frac{207}{2} + \frac{1053}{2} x \right) H_{0,0,0,0} \xi_2 - (108-156x) H_{0,0,0,0,0} \xi_2 + (108-36x) H_{4,2} - (117-75x) H_{2,0} \xi_2 \\
- (117+405x) H_{3,0,0} + (120-24x) H_{-3,-2,0} + (121+197x) H_{0,0,0,0} \left( \frac{755}{6} - \frac{11119}{6} x \right) H_{0,0,0} \right) \\
+ (135-165x) H_0 + \left( \frac{273}{2} - \frac{975}{2} x \right) H_{2,0,0,0} - (138-54x) H_{3,2} - (138+650x) H_5 \\
+ (141+39x) H_{-4,0} - \left( \frac{1421}{8} - \frac{753}{8} x \right) H_{0,4} + (216+72x) H_{-4,1,0} - (234+150x) H_{-2,-2,0} \\
- (246-174x) H_{-4,0,0} - \left( \frac{501}{2} + \frac{757}{2} x \right) H_{4,0} - (252-396x) H_{-4,2} - (260+412x) H_{2,0,0} \\
+ (272+976x) H_{-3,-2} - (276+108x) H_{-3,-1,0} + (280+632x) H_{-2,2,0} + (288+446x) H_{0,0,2} \]
\[
\left( \frac{974}{3} - \frac{2234}{3} x \right) H_0 \zeta_2 - (330 - 258 x) H_{-3,0,0,0} - (384 - 480 x) H_{-3,3} \\
+ (408 - 216 x) H_{-3,-1,0,0} - (410 + 1030 x) H_{-3,2} - (432 + 1188 x) H_{1} + (456 + 480 x) H_{-3,0,0,0} \\
- (480 + 732 x) H_{3,0} - \left( \frac{1948}{3} - \frac{2452}{3} x \right) H_3 - (684 + 516 x) H_{-2,-1,1,0} - \left( \frac{2074}{3} - \frac{2326}{3} x \right) H_{2,0} \\
+ (728 + 1432 x) H_{-2,3,0,0,0} - (757 + 1067 x) H_{-2,0,0,0,0} \\
- \left( \frac{20899}{27} - \frac{51353}{27} x \right) H_{0,0} - 108 x H_{0,0,0,0,0} - 75 x H_{0,0,0,0,0} - 48 x H_{3,0,0,0,0} + 252 x H_{0,0,0,0,0} \\
+ p_{\eta \phi}(-x) \left( -648 H_{-4,2} - 1248 H_{-3,3} + 288 H_{-3,2} + 570 H_{-2,4} - 4 H_{-2,0} + 602 H_{-2,2} \\
+ 129 H_{0,5} - 25 H_{0,4} - \frac{1037}{2} H_{0,3} + \frac{177}{4} H_{0,2} - 354 H_{0,2} \zeta_3 + 36 H_{3} - 72 H_{3,3} + 36 H_{3} \zeta_2 \\
+ 501 H_{4} - 144 H_{4} \zeta_2 + 216 H_{5} - 360 H_{6} - 318 H_{-4,0} + 648 H_{-4,2} + 1152 H_{-3,1,2} - 470 H_{-3,0} \\
- 1272 H_{-3,0,0} - 200 H_{-3,2} + 1392 H_{-3,3} + 1152 H_{-2,2} - 200 \zeta_2 + 1248 H_{-2,1,3} - 112 H_{-2,1,2} \\
+ \frac{249}{2} H_{-2,0} - 1056 H_{-2,0,0} + 316 H_{-2,0} \zeta_2 - 334 H_{-2,2} - 256 H_{-2,3} + 1752 H_{-2,4} \\
+ 1152 H_{-1,3,2} + 2548 H_{-1,3,0} - 1122 H_{-1,2,2} - 200 \zeta_2 + 570 H_{-1,0,4} - 4 H_{-1,0,3} + 602 H_{-1,0,2} \\
- 668 H_{-1,3} - 168 H_{-1,4} + 1248 H_{-1,5} - 27 H_{0,0} - 222 H_{0,0} \zeta_4 - 1550 H_{0,0,3} - 542 H_{0,0} \zeta_2 \\
+ 18 H_{0,0} - 36 H_{2,0} \zeta_3 + 18 H_{2,0} \zeta_2 - 96 H_{2,0} \zeta_2 + 167 H_{3,0} + 144 H_{3,0} \zeta_2 - 192 H_{3,1} \zeta_2 - 88 H_{3,2} \\
- 12 H_{4,0} - 264 H_{4,1} - 96 H_{4,2} - 264 H_{5,0} - 384 H_{5,1} + 324 H_{-4,0} + 176 H_{-3,1,0} \\
- 1152 H_{-3,-1,2} + 356 H_{-3,0} + 600 H_{-3,2,0} + 768 H_{-3,2,1} + 176 H_{-2,2,0} - 1152 H_{-2,2,2} \\
- 1152 H_{-2,-1,2} + 536 H_{-2,-1,0} + 1632 H_{-2,-1,0,2} + 112 H_{-2,-1,2,2} - 2112 H_{-2,-1,3} \\
- 839 H_{-2,0,0} + 1416 H_{-2,0,0} \zeta_2 - 56 H_{2,0,0} + 176 H_{2,2,1} + 384 H_{2,2,2} + 1056 H_{2,3,0} \\
+ 1152 H_{2,3,1} + 176 H_{2,-1,3,0} - 1152 H_{2,-1,2,2} + 536 H_{2,-1,2,0} \\
+ 1632 H_{2,-2,0,2} + 112 H_{2,-1,2,2} - 2112 H_{2,-1,2,3} - 1152 H_{2,-1,1,2} \\
+ 1248 H_{2,-1,1,0,2} + 224 H_{2,-1,1,3} - 2880 H_{2,-1,1,4} + 249 H_{2,-1,0,0} \\
- 864 H_{1,0,0} \zeta_3 + 344 H_{1,0,0} \zeta_2 - 334 H_{1,2,0} + 176 H_{1,2,2} + 32 H_{1,3,0} + 352 H_{1,3,1} \\
+ 384 H_{1,3,2} + 88 H_{1,4,0} + 1152 H_{1,4,1,1} - \frac{639}{4} H_{0,0,0} + 288 H_{0,0,0} \zeta_3 - 402 H_{0,0,0} \zeta_2 \\
- 96 H_{2,0} \zeta_2 + 96 H_{2,1,0} \zeta_2 - 44 H_{3,0} + 88 H_{3,1,0} - 60 H_{4,0} - 96 H_{4,1,0} - 576 H_{3,-1,0,0} \\
+ 660 H_{3,0,0,0} - 576 H_{2,-2,0,0} + 1152 H_{2,-2,1,2} + 408 H_{2,-2,2,0} - 960 H_{2,-2,2,0} \\
- 768 H_{2,-2,2,1} + 750 H_{2,-2,0,0} + 360 H_{2,2,0,0} + 384 H_{2,2,1,0} - 576 H_{2,3,0} \\
+ 1152 H_{2,-1,2,1} + 408 H_{2,-1,2,0} - 960 H_{2,-1,2,0} + 768 H_{2,-1,2,1} + 1152 H_{2,-1,1,2} \\
- 1152 H_{2,-1,1,0} \zeta_2 + 2304 H_{2,-1,1,1,3} + 1072 H_{2,-1,1,0,0} + 2112 H_{2,-1,0,0} \zeta_2 + 112 H_{2,-1,1,2,0} \\
- 768 H_{2,-1,2,2} - 1728 H_{2,-1,3,0} + 1536 H_{2,-1,3,1} + 1107 H_{2,-1,3,0,0} - 1080 H_{2,-1,0,0,0} \zeta_2 \\
+ 88 H_{2,-2,0,0} + 176 H_{2,-2,1,0} + 336 H_{2,-2,0,0} + 384 H_{2,-2,1,0} + 350 H_{0,0,0,0} + 288 H_{0,0,0,0} \zeta_2 \\
+ 36 H_{3,0,0,0} + 576 H_{2,-2,1,0} - 1200 H_{2,-1,0,0} + 744 H_{2,-2,0,0} + 576 H_{2,-2,1,0} \\
- 1200 H_{2,-1,2,0} + 576 H_{2,-1,1,2} + 1152 H_{2,-1,1,2,0} + 696 H_{2,-1,1,0,0} \\
- 768 H_{2,-1,2,0,0} - 768 H_{2,-1,2,1,0} - 816 H_{2,-1,2,0,0} - 72 H_{2,-1,2,0,0} + 510 H_{0,0,0,0,0} \\
\end{align*}
\]
\begin{align*}
+ & 1728 H_{-1,-1,-1,0,0,0} - 1344 H_{-1,-1,0,0,0,0} + 360 H_{-1,0,0,0,0,0,0} - 720 H_{0,0,0,0,0,0,0} \\
+ & p_{qg}(x) \left( 72 H_{-4} C_2 - 168 H_{-3} C_2 - 108 H_{-3} C_2 + \frac{10537}{144} H_0 - 9 H_0 C_5 - \frac{163}{2} H_0 C_4 + \frac{242}{3} H_0 C_3 \\
+ & \frac{4879}{36} H_0 C_2 + 66 H_0 C_2 C_3 - \frac{11104}{27} H_2 - 642 H_2 C_4 - 4 H_2 C_3 - 268 H_2 C_2 - \frac{3560}{9} H_3 - 480 H_3 C_3 \\
+ & 88 H_3 C_2 - \frac{947}{2} H_4 - 264 H_5 + 162 H_{-4} - 216 H_{-4,2} + 288 H_{-3} - 1 C_2 - 54 H_{-3} - 3 C_2 \\
+ & 144 H_{-3,2} - 192 H_{-3,3} + 144 H_{-2,2} - 2 C_2 - 84 H_{-2,0} C_3 - 54 H_{-2,0} C_2 + 144 H_{-2,3} - 72 H_{-2,4} \\
- & \frac{23959}{108} H_0,0 - 306 H_0,0 C_4 + \frac{1571}{3} H_0,0 C_2 + 384 H_{1,-3} C_3 - \frac{11104}{27} H_{1,0} \\
- & 642 H_{1,0} C_4 - 4 H_{1,0} C_3 + 268 H_{1,0} C_2 - 480 H_{1,2} C_3 + 348 H_{1,4} - 192 H_{1,5} + 192 H_{2,-2} C_2 \\
- & \frac{3560}{9} H_2,0 - 528 H_2,0 C_3 + 44 H_2,0 C_2 - 480 H_2,1 C_3 + 116 H_2,3 - 192 H_2,4 - \frac{477}{2} H_3,0 - 44 H_3,2 \\
- & 96 H_3,3 - 198 H_4,0 - 132 H_4,1 - 48 H_4,2 - 96 H_5,0 - 192 H_5,1 - 288 H_{-4,-1,0} + 36 H_{-4,0} \\
- & 288 H_{-3,-2,0} + 72 H_{-3,-1,0} + 96 H_{-3,1,0} + 144 H_{-3,2,0} - 192 H_{-3,3,0} - 168 H_{-3,2,0} - 192 H_{-3,3,0} \\
- & 192 H_{-2,3,0} + 36 H_{-2,2,0} - 48 H_{-2,2,0} + 144 H_{-2,1,0} - 192 H_{-2,1,0} - 1344 H_{-2,2,0} - 54 H_{-2,0,0} \\
- & 48 H_{-2,0,0} C_2 + 72 H_{-2,2,0} - 96 H_{-2,2,0} - 120 H_{-2,3,0} - 192 H_{-2,3,0} - 1733 \frac{6}{H_0,0,0} - 288 H_{0,0,0} C_3 \\
+ & 438 H_{0,0,0} C_2 - 288 H_{1,-1,0} + 288 H_{1,-3,0} - 384 H_{1,-3,2} + 192 H_{1,-2,0} C_2 - 384 H_{1,-2,3} \\
- & \frac{3560}{9} H_{1,0} - 576 H_{1,0,0} C_3 + 480 H_{1,0,0} C_3 - 288 H_{1,1,4} + 96 H_{1,2,3} + 116 H_{1,3,0} - 192 H_{1,3,2} \\
- & 336 H_{1,4,0} - 576 H_{1,4,1} - 288 H_{2,-3,0} + 144 H_{2,-2,0} - 192 H_{2,-2,2} - \frac{609}{2} H_2,0 - 96 H_{2,1,3} \\
- & 144 H_{2,3,0} - 192 H_{2,3,1} - 192 H_{3,0,0} - 72 H_{3,0,0} - 44 H_{3,1,0} + 60 H_{4,0,0} - 48 H_{4,1,0} \\
+ & 384 H_{-3,-1,-1,0} - 432 H_{-3,-1,0,0} + 84 H_{-3,0,0,0} + 192 H_{-3,-2,-1,0} - 312 H_{-2,-2,0} \\
+ & 192 H_{-2,-1,-2,0} + 72 H_{-2,-1,0,0} - 48 H_{-2,-1,0,0} - 54 H_{-2,0,0,0} - 96 H_{-2,2,0,0} - 96 H_{-2,2,1,0} \\
- & 592 H_{0,0,0,0} - 384 H_{1,-3,0,0} + 288 H_{1,-2,2,0} - 192 H_{1,-2,2,0} - \frac{1343}{2} H_{1,0,0,0} - 384 H_{1,1,1,0} \\
- & 96 H_{1,1,3,0} - 192 H_{1,2,0,0} + 60 H_{1,2,0,0} + 96 H_{1,3,0,0} - 192 H_{1,3,1,0} - 288 H_{2,-2,0,0} - 26 H_{2,0,0,0} \\
- & 192 H_{2,-2,0} + 60 H_{2,1,0,0} + 288 H_{2,2,0,0} + 300 H_{3,0,0,0} + 288 H_{3,1,0,0} + 384 H_{2,2,1,0} \\
- & 360 H_{2,-1,0,0,0} + 96 H_{2,0,0,0,0} - 510 H_{0,0,0,0,0} - 288 H_{1,-2,0,0,0} - 200 H_{1,0,0,0,0} \\
- & 384 H_{1,1,-2,0,0} + 180 H_{1,1,1,0,0} + 288 H_{1,1,2,0,0} + 624 H_{1,2,0,0,0} + 288 H_{1,2,1,0,0} + 456 H_{2,0,0,0,0} \\
+ & 624 H_{2,1,0,0,0} + 288 H_{2,1,1,0,0} + 72 H_{0,0,0,0,0,0} + 360 H_{0,0,0,0,0,0} + 768 H_{1,1,0,0,0,0} \\
+ & 864 H_{1,1,1,0,0,0} \\ & + \frac{16}{3} C_{n f}^{3} \left( (1 + x) \left( 32 H_{-2} C_2 - 60 H_{-4} - 120 H_{-2,0} + 32 H_{-1,0} C_2 + 40 H_{0,0} C_3 - 32 H_{2,2} - 56 H_{3,0} \\
- & 64 H_{3,1} + 64 H_{-2,1,0,0} + 64 H_{-1,2,0} - 240 H_{-1,0,0,0} - 368 H_{0,0,0,0} - 32 H_{2,1,0} - 16 H_{3,0} \\
+ & 128 H_{-1,-1,0,0} - 192 H_{-1,0,0,0} - 24 H_{2,0,0,0} \right) + (1 - x) \left( 32 H_{-3} C_2 + \frac{35}{2} H_0 - 20 H_2 + 32 H_5 \\
+ & 48 H_{-4,0} - 32 H_{-3,2} + 16 H_{-2,0} C_2 - 32 H_{-2,3} - 20 H_{1,0} - 64 H_{1,2} - 64 H_{2,1} + 8 H_{4,0} + 48 H_{3,0} \right) \right) \right) \right) \right) \right) \right) \right) \tag{JHEP10(2017)041} \\
- & -64 \right)
\[ -16H_{-2,2,0} - 72H_{1,0,0} - 64H_{1,1,0} + 24H_{-2,0,0,0} \] 
\[ - (4 + 28x)H_0 \xi_3 - (12 - 340x)H_{0,0,0,0} \] 
\[ - (16 + 176x)H_{-3,0} - \left( \frac{39}{2} - \frac{825x}{2} \right) H_{0,0,0,0} + (20 + 28x)H_0 \xi_4 - \left( \frac{51}{2} - \frac{29x}{2} \right) H_{0,0,0,0} \] 
\[ - (27 + 123x)H_0 \xi_2 + (32 - 48x)H_{0,0} \xi_2 - (42 - 102x)H_3 - (57 - 87x)H_{2,0} \] 
\[ - (80 + 240x)H_{-2,0,0,0} + 120xH_{0,0,0,0,0} + p_{qq}(-x) \left\{ -80H_{-2,3,0} + 184H_{-2,3}H_{0} \xi_2 - 5H_0 \xi_4 - 158H_0 \xi_3 \right. \] 
\[ + 21H_0 \xi_2 + 156H_{-4,0} - 136H_{-3,0} + 32H_{-3,2} + 64H_{-2,1,2} + 42H_{-2,0} - 16H_{-2,0} \xi_2 \] 
\[ - 104H_{-2,2} + 64H_{-2,3} + 64H_{-1,2} - 2H_{0,1,0} - 148H_{-1,0,3} + 184H_{-1,0} \xi_2 - 208H_{-1,1,3} + 96H_{-1,4} \] 
\[ + 8H_{0,0} \xi_3 - 172H_{0,0} \xi_2 + 52H_{0,3} - 32H_{3,2} - 24H_{4,0} - 96H_{4,1} + 64H_{3,0} - 144H_{-3,0} \] 
\[ + 64H_{-2,2,0} + 160H_{-2,1,0} - 64H_{-2,1,2} - 244H_{-2,0,0} + 32H_{-2,2,0} + 64H_{-2,2,1} + 64H_{-1,3,0} \] 
\[ + 160H_{-1,2,0} - 64H_{-1,2,2} + 64H_{-1,1,0} \xi_2 - 128H_{-1,1,3} + 84H_{-1,0,0} - 32H_{-1,0,0} \xi_2 \] 
\[- 104H_{-1,2,0} + 64H_{-1,2,2} + 64H_{-1,3,0} + 128H_{-1,3,1} - 63H_{0,0,0} - 48H_{0,0,0} \xi_2 - 16H_{3,0,0} \] 
\[ - 32H_{3,1} + 96H_{-2,1,1} + 8H_{-2,0,0,0} + 96H_{-1,0,0,0} - 64H_{-1,1,2} - 324H_{-1,0,0,0} + 32H_{-1,2,0,0} + 64H_{-1,2,1,0} + 112H_{0,0,0,0} + 96H_{-1,1,0,0,0} - 192H_{-1,1,0,0,0} \] 
\[ + 120H_{0,0,0,0,0} \right\} \] 
\[ + p_{qq}(-x) \left\{ - \frac{313}{8} H_0 - 15H_0 \xi_4 + 56H_0 \xi_3 + 15H_0 \xi_2 + 55H_2 - 96H_2 \xi_3 + 12H_3 \right. \] 
\[ - 120H_4 - 64H_5 + \frac{89}{2} H_{0,0} - 88H_{0,0} \xi_3 + 92H_{0,0} \xi_2 + 55H_1 + 96H_0 \xi_3 + 48H_{1,4} + 12H_{2,0} \] 
\[ - 16H_{2,0} \xi_2 + 16H_{2,3} - 40H_3 - 40H_{4,0} + 81H_{0,0,0} + 48H_{0,0,0} \xi_2 + 12H_{1,0,0} - 32H_{1,0,0} \xi_2 \] 
\[ + 16H_{1,3,0} - 52H_{2,0,0} - 112H_{0,0,0,0} - 156H_{1,0,0,0} + 32H_{1,2,2,0} + 24H_{2,0,0,0} + 32H_{2,1,0,0} \] 
\[ - 120H_{0,0,0,0,0} + 96H_{1,1,0,0,0} \right\} \] 
\[ + \frac{16}{3} C_F^2 C_A n_f \left\{ (1 + x) \left[ -16H_{-2} + 36H_{4} + 60H_{-2,0} - 16H_{-1,0} \xi_2 - 36H_0 \xi_3 + 24H_2 \right. \right. \] 
\[ + 36H_{3,0} + 48H_{3,1} - 32H_{-2,1,1} - 32H_{-1,1,0} - 32H_{-2,1,2} + 120H_{-1,0,0} + 18H_{0,0,0} \xi_2 + 24H_{2,1,0} + 12H_3 \] 
\[ - 64H_{-1,1,0,0} + 96H_{-1,1,0,0,0} + 18H_{2,0,0,0} \] 
\[ \left. + (1 - x) \left( -16H_{-3} + \frac{3161}{4} H_0 + \frac{1846}{9} H_2 \right) \right. \] 
\[ - 16H_{2} \xi_2 - 24H_{-4,0} + 16H_{-3,2} - 8H_{-2,0} \xi_2 + 16H_{-2,0} + \frac{1846}{9} H_{1,0} - 16H_{1,0} \xi_2 + 48H_{1,2} \] 
\[ + 48H_{2,1} - 24H_{-3,0,0} + 8H_{-2,2,0} + \frac{344}{3} H_{1,0,0} + 48H_{1,1,0} - 12H_{-2,0,0,0} - 12H_{1,0,0,0} \right\} \] 
\[ - (6 - 2x)H_{4,0} + (8 + 88x)H_{-3,0} - (14 + 18x)H_0 \xi_4 + (18 - 10x)H_0 \xi_3 + (20 + 28x)H_{2,0,0} \] 
\[ - (24 - 64x)H_{0,0,0} \xi_2 - (24 - 8x)H_5 + \left( \frac{77}{3} - \frac{1276}{3} x \right) H_{0,0,0,0} + (40 + 120x)H_{2,0,0} \] 
\[ - (44 + 204x)H_{0,0,0,0} - \left( \frac{145}{3} - \frac{379}{3} x \right) H_0 \xi_2 + \left( \frac{290}{3} - \frac{398}{3} x \right) H_3 + \left( \frac{317}{3} - \frac{371}{3} x \right) H_{2,0,0} \] 
\[ + (\frac{6623}{27} - \frac{8821}{27} x) H_{0,0,0,0,0} + p_{qq}(-x) \left( 40H_{-2} - 92H_{-2} \xi_2 + \frac{5}{2} H_0 \xi_4 + 79H_0 \xi_3 \right. \] 
\[ - \frac{21}{2} H_0 \xi_2 - 78H_4 + 48H_{-4,0} + 68H_{-3,0} - 16H_{-3,2} - 32H_{-2,1,2} - 21H_{-2,0} + 8H_{-2,0} \xi_2 \] 
\[ + 52H_{-2,2} - 32H_{-2,3} - 32H_{-1,2} \xi_2 + 40H_{-1,0} \xi_3 - 92H_{-1,0} \xi_2 + 104H_{-1,3} - 48H_{-1,4} \]
\[-4H_{0,0}c_3 + 86H_{0,0}c_2 - 26H_{3,0} + 16H_{3,2} + 12H_{1,0} + 48H_{4,1} - 32H_{3,-1,0} + 72H_{3,-3,0} \\
-32H_{-2,-2,0} - 80H_{-1,0} - 32H_{-2,-1,2} + 122H_{-1,2,0} - 16H_{-1,2,0,0} - 32H_{2,-2,1} - 32H_{1,-1,3} \\
-80H_{-1,-2,0} + 32H_{-1,-2,2} - 32H_{-1,1,0}c_2 + 64H_{1,1,3} - 42H_{-1,0,0} + 16H_{1,-1,0}c_2 \\
+ 52H_{1,0}c_2 - 32H_{-1,1,2} - 32H_{-1,3,0} - 64H_{-1,1,3} + \frac{63}{2}H_{0,0,0}c_3 + 8H_{3,0,0} + 16H_{3,1,0} \\
-48H_{-2,-1,0,0} + 84H_{-2,0,0,0} - 48H_{-1,-2,0,0} - 16H_{1,-1,0,0,0} + 32H_{-1,-1,2,0} + 162H_{1,0,0,0,0} \\
-16H_{1,-2,0,0,0} - 32H_{-1,1,2,0} - 56H_{0,0,0,0,0} - 48H_{-1,1,0,0,0,0} + 96H_{-1,-1,0,0,0,0} - 60H_{0,0,0,0,0,0} \right) \\
+p_{qq}(x) \left( -\frac{239}{36}H_0 + \frac{91}{2}H_0c_4 - \frac{121}{3}H_0c_3 - \frac{967}{18}H_0c_2 + \frac{2594}{27}H_2 + 88H_2c_3 - 40H_2c_2 \\
+\frac{1156}{9}H_3 - 16H_3c_2 + 95H_4 + 48H_5 + \frac{965}{27}H_0c_0 + 76H_0c_0c_3 - \frac{292}{3}H_0c_0c_2 + \frac{2594}{27}H_1c_1 \\
+88H_1c_3 - 40H_1c_2 - 24H_1c_4 + \frac{1156}{9}H_0c_0c_2 - 8H_2c_3 + 61H_3 + 8H_3c_2 + 36H_4 \\
+24H_4c_3 + \frac{769}{6}H_0c_0c_2 - 60H_0c_0c_2 + \frac{1156}{9}H_1c_1c_0 - 8H_1c_3 + 73H_2c_3 + 8H_3c_1 = 144H_0c_0,0,0 \\
+131H_{1,0,0,0} - 24H_{1,2,0,0} - 28H_{2,0,0,0} - 24H_{2,1,0,0} + 60H_{0,0,0,0,0} - 16H_{1,0,0,0,0,0} - 72H_{1,1,0,0,0,0} \right) \right) \\
+ \frac{16}{81} C_F^2 n_f^2 \left( 108(1+x)H_0c_0,0,0 + (1-x) \left( -260H_0 + 72H_0c_2 - 276H_2 - 144H_3 - 276H_{1,0} \\
- 144H_2c_0 - 144H_1c_0,0,0 \right) + \frac{159}{4}H_0 - 36H_0c_3 \\
+120H_0c_2 - 76H_2 - 240H_3 - 108H_4 + 23H_0c_0 + 72H_0c_0c_2 - 76H_{1,0} - 240H_{2,0} - 108H_{3,0} \\
-279H_{0,0,0,0} - 240H_{1,0,0,0} - 108H_{2,0,0,0} - 216H_{0,0,0,0,0} - 108H_{1,0,0,0,0,0} \right) \right). \]  
(C.1)

The most compact representation of \( \delta P^{(3)^{+}}(x) \) is via its difference to \( \delta P^{(3)^{-}}(x) \),

\[
\delta P^{(3)^{+}}(x) - \delta P^{(3)^{-}}(x) = \\
16C_F^2 (C_A - 2C_F) \left( (1+x) \left( -208H_{-2}c_3 + \frac{40}{3}H_{-2}c_2 + 192H_{-2,-1}c_2 - 376H_{-2,0} + 96H_{-2,2} \\
+ 192H_{-1,-2}c_2 - 208H_{-1,0}c_3 + \frac{40}{3}H_{-1,0}c_2 + 192H_{-1,3} + 480H_{-1,4} - \frac{400}{3}H_{2,2} - \frac{800}{3}H_{3,1} \\
+ \frac{656}{3}H_{-2,-1}c_2 - 192H_{-2,-1,2,1} + \frac{656}{3}H_{-2,-2,0} - 192H_{-2,-2,2} + 192H_{-1,-1}c_2 \\
- 384H_{-1,-1,3} - 752H_{-1,0,0} - 302H_{-1,0}c_2 + 96H_{-1,2,0} + 128H_{-1,2,2} + 288H_{-1,3,0} \\
+ 256H_{-1,3,1} - \frac{224}{3}H_{2,2,0} - \frac{400}{3}H_{2,1,0} - 96H_{-1,1,2,0} - \frac{1312}{3}H_{1,1,1}c_1 - 192H_{-1,1,2} \\
- 536H_{-1,0,0,0} + 128H_{-1,2,0,0} + 128H_{-1,2,1,0} - 248H_{-1,1,0,0,0} + 84H_{-1,0,0,0,0} \right) \\
+ (1-x) \left( 192H_{-1}c_2 + 208H_{-3}c_3 + 347H_0 - 206H_0c_5 + 64H_0c_0c_3 - \frac{452}{3}H_2 - 24H_2c_3 \\
- 60H_2c_2 + 120H_6 - 128H_{-5,0} - 144H_{-4,2} - 256H_{-3,-1}c_2 + 192H_{-3,0}c_2 - 192H_{-3,3} \\
- 128H_{-2,-2}c_2 + 104H_{-2,0}c_3 - 144H_{-2,4} - 152H_{0,0}c_4 - \frac{452}{3}H_{1,0} - 24H_{1,0}c_3 - 60H_{1,0}c_2 \right)
\]
\(- \frac{800}{3} H_{1,2} - 32 H_{1,2} \xi_2 - \frac{800}{3} H_{2,1} - 32 H_{2,1} \xi_2 + 24 H_{5,0} + 96 H_{-4,-1,0} - 184 H_{-4,0,0} \\
+ 96 H_{-3,-2,0} - 208 H_{-3,-1,0} + 192 H_{-3,-1,2} - 48 H_{-3,2,0} + 64 H_{-2,-3,0} - 104 H_{-2,-2,0} \\
+ 216 H_{-2,2,0} - 128 H_{-2,1,0} \xi_2 + 128 H_{-2,1,0} \xi_2 + 24 H_{-2,0,0} - 48 H_{-2,0,0} - 192 H_{0,0,0} \xi_3 \\
- \frac{448}{3} H_{1,0,0} - 32 H_{1,0,0} \xi_2 - \frac{800}{3} H_{1,1,0} - 32 H_{1,1,0} \xi_2 - 128 H_{-3,-1,0,0} + 24 H_{-3,-1,0,0} \\
- 216 H_{-3,0,0,0} - 64 H_{-2,-2,0,0} - 64 H_{-2,-2,0,0} + 96 H_{-2,-2,0,0} \\
- 160 H_{0,0,0,0} \xi_2 - 128 H_{-2,1,0,0} + 240 H_{-2,1,0,0} - 192 H_{2,0,0,0} + 192 H_{0,0,0,0}
\right)
\right)
+ \left(12 + \frac{3292}{3} x\right) H_{0,0,0,0} - (32 + 96 x) (H_{3,2} + H_{3,0,0} + H_{3,1,0}) - (40 - 24 x) H_{2,0} \xi_2 \\
+ (42 - 186 x) H_{0,0,0,0} - (48 - 16 x) H_{3} + \left(\frac{148}{3} - \frac{2164}{3} x\right) H_{3,0} - \left(60 - \frac{716}{3} x\right) H_{3} \\
- (96 + 288 x) H_{4,1} - \left(\frac{314}{3} - 194 x\right) H_{2,0} - \left(\frac{344}{3} + 664 x\right) H_{4,0} - \left(\frac{646}{3} + 140 x\right) H_{3,0} \\
- \left(\frac{388}{3} + 900 x\right) H_{-2,0,0} - (140 + 196 x) H_{4} + \left(\frac{1061}{6} - \frac{161}{6} x\right) H_{0,0,4} + \left(\frac{544}{3} + 416 x\right) H_{-2,2,0} \\
- \left(216 + \frac{748}{3} x\right) H_{0,0,2} + \left(\frac{704}{3} + \frac{448}{3} x\right) H_{3,3,2} - \left(\frac{800}{3} + \frac{928}{3} x\right) H_{5} + \left(279 - \frac{661}{3} x\right) H_{0} \xi_3 \\
- (304 - 112 x) H_{-2,1,0,0} - \left(\frac{988}{3} + \frac{548}{3} x\right) H_{-2,0} \xi_2 - \left(\frac{1016}{3} + \frac{136}{3} x\right) H_{3,2} \\
+ \left(\frac{1184}{3} + \frac{928}{3} x\right) H_{-2,3} + (404 - 260 x) H_{-4,0} + (408 + 72 x) H_{0,0,5} + (420 + 60 x) H_{0,0,0,0} \\
+ \left(\frac{1544}{3} + \frac{1565}{3} x\right) H_{0,0,0} + (524 + 826 x) H_{0,0,0,0} -(600 - 40 x) H_{0,0,0,0,0} + (604 - 124 x) H_{-2,0,0,0} \\
+ (608 - 320 x) H_{-3,0,0} + p_{q_0} (-x) \left(-408 H_{-4} \xi_2 - 416 H_{-3} \xi_3 + 192 H_{-3} \xi_2 + 150 H_{-2} \xi_4 \\
- \frac{304}{3} H_{-2} \xi_3 + \frac{2300}{3} H_{-2} \xi_2 + 166 H_{0} \xi_5 - \frac{335}{6} H_{0} \xi_4 - \frac{1993}{3} H_{0} \xi_3 + \frac{157}{2} H_{0} \xi_2 - 164 H_{0} \xi_3 \\
+ 24 H_{2} - 48 H_{3} \xi_3 + 24 H_{3} \xi_2 + 614 H_{4} - 72 H_{4} \xi_2 + 240 H_{5} - 320 H_{6} + 192 H_{5,0} + 460 H_{4,0} \\
+ 336 H_{4,2} + 384 H_{3,3} - 1880 H_{3,0} - 456 H_{3,0,0} \xi_2 \xi_3 \xi_4 \xi_5 \xi_6 \\
+ 384 H_{2,2,2} + 416 H_{2,1} \xi_3 + \frac{128}{3} H_{2,1} \xi_2 + 181 H_{2,0} - 448 H_{2,0} \xi_3 + \frac{616}{3} H_{2,0,0} \xi_2
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\[-384 H_{-3,1,2} - 696 H_{-3,0,0} + 240 H_{-3,2,0} + 256 H_{-3,2,1} - 48 H_{-2,3,0} + \frac{704}{3} H_{-2,2,0} \]
\[-384 H_{-2,2,2} - 384 H_{-2,1,1,0} \zeta + \frac{2144}{3} H_{-1,1,0} + 512 H_{-2,1,0} \zeta - \frac{128}{3} H_{-1,2,0} \]
\[-704 H_{-2,1,3} - \frac{3302}{3} H_{-2,0,0} - 480 H_{-2,0,1} \zeta - \frac{80}{3} H_{-2,1,1,0} + \frac{416}{3} H_{-1,2,1,0} + 128 H_{-2,2,2} \]
\[+352 H_{-2,3,1} + 384 H_{-2,1,1,0} - 144 H_{-1,1,0} - 4 \zeta + \frac{704}{3} H_{-1,1,0} + 384 H_{-1,1,0} - 384 H_{-1,2,1,0} \zeta \]
\[+ \frac{2144}{3} H_{-1,1,0} + 512 H_{-1,1,0} - 2 \zeta - \frac{128}{3} H_{-1,1,2,0} - 704 H_{-1,1,2,0} - 384 H_{-1,1,1,0} - 2 \zeta \]
\[+ 416 H_{-1,1,0} \zeta + \frac{128}{3} H_{-1,1,0} - 2 \zeta - \frac{256}{3} H_{-1,1,1,0} - 960 H_{-1,1,1,0} + 362 H_{-1,0,0} - 480 H_{-1,0,0} \zeta \]
\[+ \frac{656}{3} H_{-1,0,0} \zeta - \frac{1228}{3} H_{-1,2,0} + 16 H_{-1,2,0} \zeta + \frac{416}{3} H_{-1,1,2,0} - \frac{16}{3} H_{-1,1,3,0} + \frac{832}{3} H_{-1,1,3,1} \]
\[+ 128 H_{1,3,2} + 336 H_{1,4,0} + 384 H_{1,4,1} + \frac{507}{2} H_{0,0,0} + 384 H_{0,0,0} \zeta - 464 H_{0,0,0} \zeta \]
\[-32 H_{2,0,0} \zeta - 32 H_{2,1,0} \zeta - \frac{32}{3} H_{3,0,0} - \frac{208}{3} H_{3,1,0} - 32 H_{3,1,0} + 32 H_{1,1,0} - 288 H_{-3,1,0} \]
\[+ 432 H_{-3,0,0} - 192 H_{-2,2,0} + 384 H_{-2,1,1,0} + 448 H_{-2,1,1,0} + 300 H_{-2,1,2,0} \]
\[-256 H_{-2,2,1,0} + 128 H_{-2,1,2,0} - 288 H_{-2,1,3,0} \]
\[+ 384 H_{-2,1,2,0} + 448 H_{-1,2,0} - 320 H_{-1,2,0} + 256 H_{-1,2,2,0} + 384 H_{-1,1,2,0} \]
\[-384 H_{-1,1,1,0} - 12 \zeta + 768 H_{-1,1,1,0} - 12 \zeta + 640 H_{-1,1,1,0} + 640 H_{-1,1,0,0} - 12 \zeta - \frac{128}{3} H_{-1,1,2,0} \]
\[-256 H_{-1,1,2,0} - 576 H_{-1,1,3,0} - 1422 H_{-1,1,3,1} - 480 H_{-1,0,0} \zeta \]
\[+ \frac{64}{3} H_{-1,2,0} + \frac{416}{3} H_{-1,2,1,0} + 128 H_{-1,3,0} + 128 H_{-1,3,1} + \frac{1400}{3} H_{0,0,0} + 320 H_{0,0,0} \zeta \]
\[+ 12 H_{-1,1,0,0} - 480 H_{-2,1,0,0} + 192 H_{-1,2,1,0} - 480 H_{-1,2,2,0} \]
\[+ 192 H_{-1,1,2,0} + 384 H_{-1,1,2,0} + 640 H_{-1,1,2,0} - 256 H_{-1,1,2,2,0} - 256 H_{-1,1,2,2,0} \]
\[+ 1088 H_{-1,0,0,0} + 740 H_{0,0,0,0} + 576 H_{-1,1,0,0} - 768 H_{-1,1,0,0} + 560 H_{-1,0,0,0} \]
\[-288 H_{0,0,0,0} \zeta \]
\[+ \frac{16}{3} C_F^2 (C_A - 2 C_F)^2 \left\{ (1+x) \left( -1248 H_{-2,2,0} - 316 H_{-2,2,1} + 1152 H_{-2,1,2,0} - 834 H_{-2,0,0} \right) \right. \]
\[+ 552 H_{-2,2,2} + 1152 H_{-1,2,1,0} - 1248 H_{-1,1,1,0} - 316 H_{-1,1,1,0} \zeta + 1104 H_{-1,1,3,0} + 2880 H_{-1,1,3,0} \]
\[-560 H_{2,2} - 1120 H_{3,1} + 472 H_{2,1,1} - 1152 H_{2,1,2} + 768 H_{2,2,1} + 1247 H_{1,2,0} \]
\[-1152 H_{1,2,2} + 1152 H_{1,2,1} + 2034 H_{-1,1,2,0} - 1168 H_{-1,1,2,0} - 2112 H_{-1,1,2,0} \zeta \]
\[+ 552 H_{-1,2,2} + 768 H_{1,2,2} + 1728 H_{1,3,3} + 5136 H_{1,3,1,0} - 560 H_{1,1,2,0} - 576 H_{1,1,2,0} \]
\[+ 944 H_{1,1,0,0} - 1152 H_{1,1,1,0} - 408 H_{1,1,2,0} + 768 H_{1,2,0} + 768 H_{1,2,0,0} + 72 H_{2,0,0} \]
\[-1728 H_{1,0,0,0} + 1344 H_{1,0,0,0} \bigg) \left( 1 + x \right) \left( 1 + \frac{2445}{2} H_{0,0,0} \right) \]
\[-126 H_{0,0,0} + 264 H_{0,1,1,0} - 766 H_{2,2,0} - 156 H_{2,1,2} - 156 H_{2,1,0} + 14 H_{1,2,1} + 300 H_{0,0,0} \]
\[-48 H_{-5,0} + 390 H_{-4,0} - 568 H_{-4,2} + 960 H_{-3,1} - 1248 H_{-3,0,0} - 816 H_{-3,0,0} \zeta - 864 H_{-3,3} - 480 H_{-2,2,0} \]
\[+ 372 H_{-2,2,0} - 484 H_{-2,1,1,0} - 33 H_{0,0,0} - 766 H_{1,0} + 72 H_{1,0} - 156 H_{1,0} - 1120 H_{1,2} \]
\[
-192H_{1,2}\zeta_2 - 192H_{2,0}\zeta_2 - 1120H_{2,1} - 192H_{2,1,0}\zeta_2 + 48H_{3,0}\zeta_2 + 60H_{5,0} + 144H_{-4,-1,0}
-420H_{-4,0,0} + 144H_{-3,-2,0} - 168H_{-3,-1,0} + 864H_{-3,-1,2} - 216H_{-3,2,0} + 96H_{-2,-3,0}
-84H_{-2,-2,0} + 432H_{-2,-2,2} - 480H_{-2,-1,0}\zeta_2 + 864H_{-2,-1,3} + 576H_{-2,0,0}\zeta_2 - 216H_{-2,3,0}
-252H_{0,0,0}\zeta_3 - 944H_{1,0,0} - 192H_{1,0,0,0}\zeta_2 - 1120H_{1,1,0} - 192H_{1,1,0,0}\zeta_2 - 192H_{-3,-1,0}
+624H_{-3,-1,0,0} - 588H_{-3,0,0,0} - 96H_{-2,-2,0} + 360H_{-2,-2,0,0} - 96H_{-2,-1,0,0}
+432H_{-2,-1,2,0} - 264H_{0,0,0,0}\zeta_2 + 144H_{1,0,0,0} - 192H_{-2,-1,0,0} + 792H_{-2,-1,0,0}
-384H_{-2,0,0,0,0} + 108H_{0,0,0,0,0,0} - 144 + 528xH_{3,0,0} + (162 + 1158x)H_{0,0}\zeta_2
-(174 + 956x)H_{0}\zeta_2 - (192 + 576x)(H_{3,2} + H_{3,1,0}) - (196 + 1100x)H_{-3,0}
-(799/4 + 3761/4x)H_{0}\zeta_4 + (324 + 428x)H_{0,0,0,0} - (332 + 1236x)H_{-2,0,0} - (424 + 520x)H_{2,0,0}
-(424 + 1352x)H_{4,0} + (447 + 469x)H_{0}\zeta_3 - (544 + 1952x)H_{5} - (552 + 2256x)H_{4}
+(576 + 1152x)H_{0,0}\zeta_3 - (576 + 1728x)H_{4,1} - (645 + 75x)H_{0,0,0,0} - (648 - 1240x)H_{3}
-(656 + 1224x)H_{3,0} + (657 + 1503x)H_{0,0,0}\zeta_2 - (744 + 408x)H_{-2,-1,0} - (796 - 1092x)H_{2,0}
+(800 + 1120x)H_{-2,2,0} + (832 + 1472x)H_{3,-3,2} - (916 + 1388x)H_{-3,0,0,0} - (1209 + 1869x)H_{0,0,0,0,0} + (1350 + 1050x)H_{-2,0,0,0} + (1486 + 2578x)H_{0,0}
-(1514 + 1750x)H_{-2,0,0,0} + (1792 + 2432x)H_{-2,3,1,0} - p_{qq}(x) - 1296H_{-1,2,0} - 2496H_{-3,3,0}
+576H_{-3,2,0} + 1140H_{-2,3,0} - 8H_{-2,2,0} + 1204H_{-2,2,0} + 258H_{0,0}\zeta_5 - 50H_{0}\zeta_4 - 1037H_{0}\zeta_3
+(177/2)H_{0}\zeta_2 - 708H_{0,0}\zeta_3 + 72H_{3} - 144H_{3,2} + 72H_{3,2} + 1002H_{4} - 288H_{4,2} + 432H_{5}
-720H_{5} - 636H_{-4,0} + 1296H_{-4,2} + 2304H_{-3,2,0} - 940H_{-3,3,0} - 2544H_{-3,0}\zeta_2 - 400H_{-3,2}
+2784H_{-3,3,0} + 2304H_{-2,0,0} + 2496H_{-2,0,0,0} - 224H_{-2,1,0,0} - 249H_{-2,0,0,0} - 2112H_{-2,0,0,0}
+632H_{-2,0,0,0} - 668H_{-2,2,0} - 512H_{-2,3,0} + 3504H_{-2,4,0} + 2304H_{-1,3,2,0} + 2496H_{-1,2,0,0}
-224H_{-1,2,0,0} + 1140H_{-1,0,0,0} - 8H_{-1,0,0,0} + 1204H_{-1,0,0,0} - 1336H_{-1,0,0,0}
+336H_{-1,0,0,0} + 2496H_{-1,0,0,0} - 54H_{0,0,0,0} - 444H_{0,0,0,0} - 316H_{0,0,0,0} - 1084H_{0,0,0,0} + 36H_{2,0,0,0}
+36H_{2,0,0,0} - 192H_{2,2,0,0} + 334H_{3,0,0} - 288H_{3,0,0} - 38H_{3,1,0,0} - 176H_{3,2,0} - 24H_{4,0,0} - 528H_{4,1}
-192H_{4,2,0} - 528H_{5,0} - 768H_{5,1} + 648H_{-4,0} + 352H_{-3,1,0} - 2304H_{-3,2,0} - 1152H_{-3,3,0,0}
+1200H_{-3,0,0} + 1536H_{-3,1,0} + 352H_{-2,0,0} - 2304H_{-2,0,0} - 2304H_{-2,0,0,0} - 2304H_{-2,1,0,0}
+1072H_{-2,1,0,0} + 3264H_{-2,1,0,0} + 224H_{-2,1,0,0} - 128H_{-2,1,0,0} - 1768H_{-2,0,0,0}
-2832H_{-2,0,0,0} - 112H_{-2,2,0} + 352H_{-2,2,0} + 768H_{-2,2,0} + 2112H_{-2,3,0} + 2304H_{-2,3,0}
-192H_{1,0,0,0} - 304H_{1,0,0,0} - 2304H_{1,0,0,0} - 128H_{1,0,0,0} - 3264H_{1,0,0,0} - 304H_{1,0,0,0}
+224H_{1,0,0,0} - 2424H_{-1,2,3} - 2304H_{-1,2,3} - 2304H_{-1,2,3,0} - 2496H_{-1,1,0,0} - 224H_{-1,0,0,0}
+448H_{-1,1,0,0} - 576H_{1,1,0,0} + 498H_{-1,0,0,0} - 1728H_{-1,0,0,0} + 688H_{-1,0,0,0} - 688H_{-1,0,0,0}
+352H_{-1,2,0} + 64H_{-1,3,0} + 704H_{-1,3,1,0} + 768H_{-1,3,2,0} + 1776H_{-1,4,0} + 2304H_{-1,4,0}
-639H_{0,0,0,0} + 576H_{0,0,0,0} - 804H_{0,0,0,0} - 192H_{2,0,0,0} - 192H_{2,1,0,0} - 88H_{3,0,0} - 176H_{3,1,0}
-120H_{4,0,0} - 192H_{4,1,0,0} - 1152H_{-3,1,0} - 1320H_{-3,0,0,0} - 1152H_{-2,2,0,0} + 2304H_{-2,1,0,0} - 1,2
A difference between the time-like and space-like case appears for the quantities $P_{ns}^{(3)_s}$ for the first time at the four-loop level with

$$\delta P^{(3)_s}(x) = -\frac{16}{3} n_f C_F \frac{d^{abc} d_{abc}}{N_R} \left\{ (1+x) \left[ -336H_{-2} - 2\zeta_5 + 648H_{-2} \zeta_2 + 384H_3 \zeta_3 + 576H_4 \zeta_2 - 328H_{-2,0},0 \right] \right\}.$$
\[ -256H_{-2,-2}-336H_{-1,0} \zeta_3 + 648H_{-1,0} \zeta_2 - 512H_{-1,1} + 1440H_{-1,4} + 192H_{2,0} \zeta_3 + 1312H_{2,2} - 384H_{2,2} \zeta_2 - 288H_{2,4} - 384H_{3,0} \zeta_2 + 2624H_{3,1} - 768H_{3,1} \zeta_2 - 192H_{3,3} + 576H_{4,2} + 2304H_{5,1} + 784H_{-2,-1}-192H_{-2,-1,2} + 786H_{-2,-2,1} + 384H_{-1,-3,0} + 784H_{-1,-2,0} - 192H_{-1,-2,2} - 384H_{-1,-1,-3,0} - 768H_{-1,0,0} \zeta_2 - 256H_{-1,2,0} + 768H_{-1,2,2} + 1248H_{-1,3,0} + 1536H_{-1,3,1} - 192H_{2,0,0} \zeta_2 + 1312H_{2,1,0} - 384H_{2,1,0} \zeta_2 - 96H_{2,3,0} + 432H_{4,0,0} + 576H_{4,1,0} - 384H_{-2,-1,-1,0} - 384H_{-1,-2,-1,0} + 480H_{-1,-2,0,0} - 384H_{-1,-1,-2,0} + 1568H_{-1,-1,0,0} - 192H_{-1,-2,0} - 1560H_{-1,0,0,0} + 768H_{-1,2,0,0} + 768H_{-1,2,1,0} - 96H_{2,2,0,0} - 336H_{3,0,0,0} - 192H_{3,1,0,0} - 768H_{-1,-1,0,0} + 384H_{-1,0,0,0,0} - 384H_{2,0,0,0,0} - 288H_{2,1,0,0,0} + (1 - x) \left( 112H_{-4,2} + 1344H_{-3,3} \zeta_3 + 332H_{0} + 468H_{2} + 288H_{2,3} - 392H_{2,2} - 384H_{5,0} - 576H_{4,2} - 1536H_{3,3} - 1152H_{3,0} \zeta_2 - 768H_{3,3} - 768H_{2,-2,0} + 672H_{2,0} \zeta_3 - 576H_{2,-2,1} + 468H_{1,0} + 288H_{0,0} \zeta_3 - 392H_{1,0} \zeta_2 + 2912H_{1,2} - 432H_{1,4} + 2912H_{2,1} - 144H_{2,3} + 1152H_{2,4,-1,0} - 1056H_{4,0,0} + 1152H_{3,-2,0} + 768H_{3,-2,0,0} + 768H_{2,-2,0,0} + 768H_{2,-1,3} + 768H_{2,0,0} \zeta_2 - 192H_{2,3,0} + 2912H_{1,0,0} + 2912H_{1,1,0} - 144H_{1,3,0} - 1536H_{3,-1,1,0} + 1920H_{-3,-1,0,0} - 1248H_{-3,0,0,0} - 768H_{-2,-2,1,0} + 1344H_{-2,-2,0,0} - 768H_{-2,-1,2,0} + 384H_{-2,-1,0,0} - 396H_{1,0,0,0} - 144H_{1,2,0,0} - 144H_{2,1,0,0} - 1536H_{3,-2,-1,-1,0} + 1728H_{-2,-1,0,0,0} - 768H_{-2,0,0,0,0} - 576H_{1,0,0,0,0} - 432H_{1,1,0,0,0,0} \right) - (96 - 1056x + 1024x^2)H_{-2,-1,0,0} + (96 - 864x - 512x^2)H_{3,2} + H_{3,3} + 96x + 512x^2)H_{1,0,0} \zeta_2 + (96 + 192x + 512x^2)H_{2,0} \zeta_2 + (96 + 480x + 512x^2)H_{3} \zeta_2 + (96 + 672x - 512x^2)H_{2,-2,0} - (144 - 240x)H_{0} \zeta_5 - (192 - 960x + 1024x^2)H_{-3,-1,0} - (192 - 960x - 1536x^2)H_{4,1} + (288 - 576x + 678x^2)H_{-1,0} - (288 - 432x)H_{2,0,0} + (372 - 1680x - 736x^2)H_{0} \zeta_4 - (400 + 1168x)H_{-3,0} - (480 + 1280x^2)H_{0,0,0,0} + (576 - 960x + 1536x^2)H_{3,0,0} + (576 - 192x + 1024x^2)H_{3,2} - (576 + 1152x + 2048x^2)H_{5} - (976 + 1504x)H_{0} \zeta_3 + (672 - 672x + 1536x^2)H_{-2,0,0} - (672 - 672x + 1536x^2)H_{-3,0} - (1184 + 1856x)H_{0,0,0,0} + (672 + 288x + 512x^2)H_{2,2,0} - (720 + 48x + 768x^2)H_{-2,0} \zeta_2 - (2280 + 1256x)H_{0,0} \zeta_2 + (720 + 864x + 2304x^2)H_{0,0,0,0} + (768 + 1792x^2)H_{0,0,0,0} - (920 + 1688x)H_{-2,0,0} + (1056 + 288x + 1024x^2)H_{2,3,0} + (1108 + 1444x)H_{2,0,0} + (1488 + 1296x)H_{5,0} + (1600 + 720x)H_{6} - (1728 + 576x)H_{0,0,0,0} \zeta_2 + (2252 + 1940x)H_{3,0} + (2280 + 1524x)H_{0,0,0,0} - (2376 + 2040x)H_{0,0,0,0} - (2724 + 516x)H_{0,0,0} + (2820 + 1884x)H_{4} + (3232 + 6136x)H_{0,0,0} + (3582 - 1972x)H_{2,0} + (4792 - 1032x)H_{3} - (1584 + 1280x^2)H_{4,1} - (672 + 512x^2)H_{3,0,0} + \left( x + x^2 \right) \left( -896H_{-2} \zeta_3 + 768H_{-2,-1,0} \zeta_2 - 640H_{-2,0} \zeta_2 + 512H_{-2,3} + 768H_{-1,2} \zeta_2 - 896H_{-1,0} \zeta_3 + 1536H_{-1,4} - 256H_{2,1} \zeta_2 - 128H_{2,0} \zeta_2 - 256H_{2,1} \zeta_2 - 256H_{-2,0} - 512H_{-2,-1,2} + 256H_{-2,2,0} + 512H_{-2,2,1} - 512H_{-1,-3,0} - 512H_{-1,-2,2} \right) \]
\begin{align*}
+ 768 H_{-1,-1,0} \zeta_2 & - 1024 H_{-1,-1,3} - 1280 H_{-1,0,0} \zeta_2 + 512 H_{-1,2,2} + 1024 H_{-1,3,0} \\
+ 1024 H_{-1,3,1} - 256 H_{1,0,0} \zeta_2 - 256 H_{1,1,0} \zeta_2 + 512 H_{-2,-1,1,0} - 512 H_{-2,-1,0,0} \\
+ 256 H_{-2,0,0,0} + 512 H_{-1,-2,1,0} - 1024 H_{-1,-2,0,0} + 512 H_{-1,-1,2,0} - 512 H_{-1,-1,1,2,0} \\
+ 512 H_{-1,2,0,0} + 512 H_{-1,2,1,0} + 1024 H_{-1,1,1,0,0} - 1536 H_{-1,1,0,0,0} + 1024 H_{-1,0,0,0,0} \\
- 1152 H_{0,0,0,0} \zeta_3 + 576 H_{0,0,0,0,0,0}
\end{align*}

\begin{equation}
(C.3)
\end{equation}

\section{The complete $\zeta_5$ contributions}

Here we finally present the exact expressions for the part of $\gamma^{(3)\pm}_{\text{ns}}(N)$ which is proportional to $\zeta_5$:

\begin{align*}
\gamma^{(3)+}_{\text{ns}}(N) \bigg|_{\zeta_5} &= C_F^4 320 \zeta_5 \left( \frac{111}{12} + 6 \eta - 9 \eta^2 + 14 S_{-2} \right) \\
&+ C_F^3 C_A 320 \zeta_5 \left( - \frac{59}{4} - 12 \eta + 18 \eta^2 - 22 S_{-2} \right) \\
&+ C_F^2 C_A^2 80 \zeta_5 \left( \frac{67}{3} + \frac{67}{3} \eta + \frac{58}{3} S_1 - 58 \eta^2 + 8 S_1 \eta - 8 (S_1)^2 + 34 S_{-2} \right) \\
&+ C_F C_A^3 80 \zeta_5 \left( \frac{13}{4} + \frac{40}{3} \eta - \frac{116}{3} S_1 + 43 \eta^2 - 16 S_1 \eta + 16 (S_1)^2 - 6 S_{-2} \right) \\
&+ \frac{d_F^R d_F^A}{N_R} 320 \zeta_5 \left( - \frac{25}{2} + \frac{25}{3} \eta + \frac{58}{3} S_1 - 23 \eta^2 + 8 S_1 \eta - 8 (S_1)^2 \right) \\
&+ n_f C_F^3 160 \zeta_5 \left( \frac{3}{2} + \eta - 2 S_1 \right) + n_f C_F^2 C_A 80 \zeta_5 \left( - \frac{3}{2} - \eta + 2 S_1 \right) \\
&+ n_f C_F C_A^2 80 \zeta_5 \left( - \frac{33}{2} + 25 \eta + 26 S_1 + 12 \eta^2 \right) \\
&+ n_f \frac{d_F^R d_F^A}{N_R} 1280 \zeta_5 \left( 3 - 5 \eta - 2 S_1 + 6 \eta^2 \right), \tag{D.1}
\end{align*}

\begin{align*}
\gamma^{(3)-}_{\text{ns}}(N) \bigg|_{\zeta_5} &= \gamma^{(3)-}_{\text{ns}}(N) \bigg|_{\zeta_5} \\
&+ C_F^4 320 \zeta_5 \left( - \frac{13}{6} + \frac{29}{6} \eta + 7 \eta^2 \right) + C_F^3 C_A 320 \zeta_5 \left( \frac{1}{3} - \frac{20}{3} \eta - 11 \eta^2 \right) \\
&+ C_F^2 C_A^2 80 \zeta_5 \left( - 1 + 4 \eta + 17 \eta^2 \right) + C_F C_A^3 80 \zeta_5 \left( \frac{5}{6} + \frac{47}{6} \eta - 3 \eta^2 \right) \\
&+ \frac{d_F^R d_F^A}{N_R} 320 \zeta_5 \left( - \frac{1}{6} - \frac{13}{6} \eta \right). \tag{D.2}
\end{align*}

Since only sums with $w \leq 2$ can occur with $\zeta_5$, the corresponding functional form is so restricted that a direct determination and verification is possible with eight even and odd moments.
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