Evidence for gapped spin-wave excitations in the frustrated Gd$_2$Sn$_2$O$_7$ pyrochlore antiferromagnet from low-temperature specific heat measurements

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We have measured the low-temperature specific heat of the geometrically frustrated pyrochlore Heisenberg antiferromagnet Gd$_2$Sn$_2$O$_7$ in zero magnetic field. The specific heat is found to drop exponentially below approximately 350 mK. This provides evidence for a gapped spin-wave spectrum due to an anisotropy resulting from single ion effects and long-range dipolar interactions. The data are well fitted by linear spin-wave theory, ruling out unconventional low-energy magnetic excitations in this system, and allowing a determination of the pertinent exchange interactions in this material.

In the insulating R$_2$M$_2$O$_7$ magnetic pyrochlores, the rare-earth ions (R$^{3+}$) sit on a lattice of corner-sharing tetrahedra, resulting in geometrical frustration for antiferromagnetic nearest-neighbor exchange between Heisenberg spins. As such, these materials display a wide variety of interesting and unusual behaviors. Examples of observed phenomenology include spin ice [1], spin glass [2, 3], and spin liquid [4] behaviors as well as long-range magnetic order [5, 6, 7, 8, 9, 10, 11] and perhaps some novel type of hidden order [12]. A common trend observed in all of these systems is an apparent persistence of spin dynamics down to the lowest temperatures as revealed by muon spin relaxation ($\mu$SR) [3, 4, 7, 8, 9, 10, 11] and neutron spin echo [1] experiments. The origin of persistent spin dynamics (PSDs) is a major open question in the study of highly frustrated insulating magnetic oxide materials [3].

Perhaps what is most perplexing, is that PSDs have been found in pyrochlores, such as Gd$_2$Ti$_2$O$_7$ (GTO) [9, 11] and Gd$_2$Sn$_2$O$_7$ (GSO) [7, 8, 13], which, according to neutron scattering experiments [6, 10], display long-range magnetic order below a critical temperature $T_c = 0.74$ K for GTO [4] and $T_c = 1.0$ K for GSO [10]. This is highly unusual since conventional wisdom states that collective magnon-like excitations, hence spin dynamics, should freeze out in the limit of zero temperature. Similarly, the Gd$_3$Ga$_5$O$_{12}$ antiferromagnetic garnet (GGG), which displays rather extended magnetic correlations below 140 mK [14] that are now theoretically rationalized [15], also exhibits PSDs below 100 mK [16]. It is tempting to speculate that the unusual low-energy excitations giving rise to PSDs may result from a remnant of the high frustration that these systems possess and the lack of long-range order that would occur were it not for perturbatively small exchange interactions beyond nearest-neighbor [17], dipolar interactions [18, 19] or single-ion anisotropy [20].

We see Gd$_2$Sn$_2$O$_7$ (GSO) as a key system to investigate in seeking to convincingly demonstrate, via measurements of the low temperature magnetic specific heat, $C_m(T)$, the existence of unconventional low-energy excitations, as suggested by the observation of PSDs. Previous $C_m$ measurements on this material between 350 mK and 800 mK (see Fig. 1) were found to be parametrized by a $C_m(T) \sim T^2$ law [3]. Such a temperature dependence of $C_m$ is unusual since conventional antiferromagnetic magnon excitations without a gap lead to $C_m(T) \sim T^3$ or, with an anisotropy energy gap $\Delta$, to $C_m(T) \sim \exp(-\Delta/T)$. Such an unconventional $C_m(T)$ behavior further argues for the existence of unusual low-energy excitations in GSO. Interestingly, the related GTO, which also exhibits PSDs, displays $C_m(T) \sim T^2$ for 100 mK $\lesssim T \lesssim 500$ mK [3].

In contrast, a recent theoretical work [19] argues that the lowest temperature $\sim 350$ mK considered in Ref. 3 for GSO corresponds to the upper temperature limit above which multi-magnon excitations start to proliferate, obscuring the genuine single-magnon/low-temperature regime. In particular, calculations using the microscopic spin Hamiltonian 10, 13, 19 that describes the experimentally observed ground state of GSO 10 predict that $C_m$ should begin to drop away exponentially exactly at or just below 350 mK as gapped magnetic excitations become quenched at lower temperatures 19.

In this letter we present an investigation of the nature of the low-energy spin excitations in GSO through measurements of $C_m(T)$ down to 115 mK. The results show an exponentially dropping specific heat at low temperature and, when compared with theoretical calculations 19, confirm the picture of conventional gapped antiferromagnetic magnons in Gd$_2$Sn$_2$O$_7$.

Polycrystalline samples of Gd$_2$Sn$_2$O$_7$ were prepared by solid state reaction as in Ref. 21. Gd$_2$O$_3$ and SnO$_2$...
were mixed in stoichiometric ratio, pressed into pellets, and then fired in air for 48 hours at 1350°C. The pellets were then reground and the process repeated. Powder X-ray diffraction spectra taken after the two firings were indistinguishable. All peaks were indexed to space group Fd3m, with the size a of the cubic unit cell a = 10.460 Å, with no impurity peaks detectable at the 1% level. High temperature susceptibility measurements yielded a magnetic moment of 7.95μB for Gd³⁺ in GSO, close to the expected free-ion moment of 7.94μB.

Specific heat measurements were performed using the quasi-adiabatic technique. A 1 kΩ RuO₂ resistor was used as thermometer and calibrated to a commercially calibrated germanium resistance thermometer. A 10 kΩ metal-film resistor was used as a heater. The sample was suspended from thin nylon threads and the heater and thermometer were fixed directly to the sample. Leads to the thermometer and heater were made from 6 μm diameter, 1 cm long Nb/Ti wires. The sample was weakly heat-sunk to the mixing chamber of a ³He/⁴He dilution refrigerator with Pt₀.₉₂W₀.₀₈ wire, chosen for its insignificant contribution to the addendum. The addendum, due to other components such as heater and thermometer, was determined to be less than 2% of the sample’s heat capacity in the worst case at around 120 mK and thus does not effect the results of our analysis.

Slow thermal relaxation within the sample was observed (with a relaxation time constant τ₁ ~ 120 s at 200 mK). The thermal link was chosen so that the relaxation time of the sample temperature to cryostat temperature, τ₂, would be much longer than τ₁, thereby minimizing the temperature gradient within the sample. A double exponential form was fit to the sample’s temperature as a function of time after a heat pulse and the longer exponential was extrapolated to zero time, giving a measure of the sample’s heat capacity. The data presented here was taken using a long time constant τ₂ ≈ 20τ₁. Another experiment performed with τ₂ ≈ 5τ₁ resulted in a slightly noisier measurement. Yet, the specific heat in both measurements overlapped within the estimated error bars of the data, ruling out significant inaccuracies due to the slow thermal relaxation in the sample.

Our specific heat results, shown in Fig. 1, are largely consistent with previous work [5], though there is a 10 to 15% systematic discrepancy between the data sets. Our measurements, however, extend to considerably lower temperature (~ 115 mK) allowing us to test the proposal of gapped magnon excitations [19]. We observe below ~ 350 mK a deviation from the T² behavior describing the data between 350 mK and 800 mK as previously reported [5]. The specific heat decreases also faster than the T³ power-law expected for conventional gapless antiferromagnetic magnons. This already suggests that the specific heat may be dropping out exponentially, indicating a gapped spin-wave spectrum.

At the lowest temperatures an upturn in C(T) becomes apparent and can be ascribed to a nuclear contribution, C_N(T), stemming largely from the nuclear electric quadrupole interaction and, to a much lesser extent, to the nuclear hyperfine interaction. By properly parametrizing this term and subtracting it from the total specific heat, we can isolate the contribution from the electronic moments. ¹⁵⁵Gd Mössbauer spectroscopy experiments find a nuclear quadrupole electric splitting of -4.0 mm s⁻¹ [13]. Moreover, the four nuclear states are split with energies 0, 0.05, 12.1 and 15.9 mK as a result of both nuclear quadrupole and nuclear hyperfine interactions [13]. The only other naturally occurring isotope that contributes to the nuclear specific heat is ¹⁵⁷Gd which, like ¹⁵⁵Gd, has spin I = 3/2. The ratio of the nuclear dipole moments is μ¹⁵⁷_N / μ¹⁵⁵_N ≃ 0.8 and the ratio of quadrupole moments is Q¹⁵⁷_N / Q¹⁵⁵_N ≃ 1.1 [22]. The quadrupole interaction is expected to be proportional to Q_N I_z² where the local [111] z-direction points to the center of the Gd tetrahedra. There are no I_x and I_y terms in the nuclear electric quadrupole interaction due to axial symmetry of the Gd magnetic site [22]. The nuclear hyperfine interaction is proportional to μ_N I · J. However, in the Palmer-Chalker ground state of GSO [19], the electronic moments are ordered perpendicular to the local z-direction, so we can simply write the nuclear hyperfine term proportional to μ_N I_z. Thus the total nuclear Hamiltonian is H_N = aμ_N I_z + bQ_N I_z². We have calculated the nuclear contribution, C_N(T), to C(T) by choosing coefficients a and b such that the resulting eigenvalues of H_N match the above energy splittings for ¹⁵⁵Gd.

FIG. 1: (Color Online) Specific heat of Gd₂Sn₂O₇ as a function of temperature from this work (blue circles) between 115 and 800 mK. Previously measured data at higher temperatures (green diamonds) [5], the T² power law previously proposed (dotted red line) and a T⁴ power law (dashed blue line) are also plotted. The upturn seen below 150 mK results from the nuclear electric quadrupole interaction (solid line).
The interactions are then scaled by the ratios of nuclear moments to get the contribution from $^{157}$Gd. The calculated specific heats are multiplied by the relative natural isotopic abundances and added to give the total $C_N(T)$.

At temperatures much higher than the nuclear energy splittings, the nuclear contribution behaves as $C_N = AT^{-2}$ with, based on the calculation above, $A \approx 1.35 \times 10^{-4}$ JK/mol Gd. However, because of uncertainties in the nuclear moments, we keep $A$ as a free parameter. $C_N(T)$ can be subtracted from the total $C(T)$ at these temperatures since the hyperfine coupling is only a few mK. As a first rough indicator of a gapped excitation spectrum, we note that the electronic part of $C$ can be parametrized at low temperature by $C_e(T) \propto \beta T^\alpha$ as in conventional colinear antiferromagnets with single ion anisotropy. A plot of $\log(CT^2)$ versus $1/T$ should then give a straight line and this appears to be the case if $A$ is chosen to be $1.63 \times 10^{-4}$ JK/mol Gd as shown in Fig. 2 which is close to the value calculated above. The discrepancy is reasonable considering the uncertainty on the parameters in the calculation and the limited amount of data in the temperature range where the nuclear contribution becomes dominant. The phenomenological “gap” $\Delta$ so obtained is $\Delta = 1.40 \pm 0.01$ K. We now go on to describe a microscopic calculation which further supports the conclusion that a gapped magnon excitation spectrum exists in GSO.

$\text{Gd}_2\text{Sn}_2\text{O}_7$, and the related material $\text{Gd}_2\text{Ti}_2\text{O}_7$, are well approximated as isotropic, Heisenberg antiferromagnets as the Gd$^{3+}$ ions have half-filled shells ($L = 0$, $S = 7/2$). Neutron scattering experiments have identified a $k = (0, 0, 0)$ ordered ground state [10] – the so-called Palmer-Chalker state [18]. Thus $\text{Gd}_2\text{Sn}_2\text{O}_7$ is particularly well suited to the standard protocol of, after having identified the ground state, computing the second-quantized (magnon) low-lying spin excitations around that state [19, 24]. The nearest neighbor exchange interaction is estimated from the Curie-Weiss constant $\theta_{\text{CW}} = -8.6$ K by $J_1 \approx 3\theta_{\text{CW}}/z(S+1) = -0.273$ K where $z = 6$ is the number of nearest neighbors. The strength of the dipolar interaction is derived from the size of the magnetic moments and the geometry of the lattice [10, 18, 19, 24]. Though the moments are fairly isotropic, there exists a small crystal field anisotropy resulting from second order admixing from the electronic $L=0$ state to excited manifolds [25]). The crystal field parameter $B_0^2 = 47$ mK was used in this calculation as discussed in Ref. [19].

A general two-body spin interaction Hamiltonian which includes isotropic magnetic exchange interactions up to third nearest neighbor and anisotropy in the form of interactions with the local crystal field as well as long range dipole-dipole interactions [10, 18, 19, 24] was considered. The techniques of linear spin-wave theory [19, 24] were employed to diagonalize the Hamiltonian and arrive at a Bose gas of non-interacting magnons created (annihilated) by operators $a_k^\dagger (a_k)$,

$$\mathcal{H} = \mathcal{H}^{(0)} + \sum_k \sum_a \epsilon_a(k) \left[a_k^\dagger (a_k) + \frac{1}{2}\right], \quad (1)$$

where $\mathcal{H}^{(0)}$ is the classical ground state energy and $\epsilon_a(k)$ are the spin-wave dispersion energies. The pyrochlore lattice can be viewed as a FCC lattice with a 4-site ba-
The peculiar $T^2$ behavior of the specific heat of GTO also remains a mystery for, in contrast to GSO, it does not exhibit a gap down to 100 mK as might be naively expected. It would be valuable to investigate the possibility that the $T^2$ power law in GTO also gives way to an exponentially decaying specific heat at lower temperature. GSO and GTO are perhaps at this time the two pyrochlore materials most amenable to matching theory with experiment. Hence determining the cause of PSDs in these pivotal systems may provide insight into the cause of persistent spin dynamics in other more exotic, highly frustrated systems such as Tb$_2$Tb$_2$O$_7$, Yb$_2$Ti$_2$O$_7$, and Gd$_3$Ga$_5$O$_{12}$. We thank Rob Kiefl for stimulating discussions. Funding for this research was provided by NSERC, CRC (Tier 1, M.G.), CFI, MMO and Research Corporation. We thank P. Bonville for kindly providing us with specific heat data for the purpose of comparison. M.G. acknowledges the U. of Canterbury (UC) for support and the hospitality of the Dept. of Physics and Astronomy at UC where part of this work was completed.

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