First principles investigation on anomalous lattice shrinkage of W alloyed rock salt GeTe

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ABSTRACT

According to Vegard’s law, larger radius atoms substitute for smaller atoms in a solid solution would enlarge the lattice parameters. However, by first-principles calculations, we have observed unusual lattice shrinkage when W replaces Ge in rock salt GeTe. We attribute this anomalous contract to the larger electronegativity difference between W and Te than between Ge and Te, which results in shorter W-Te bonds and pronounced local distortion around W dopants. The present work would provide new insight into the lattice parameter determination and a deeper understanding of the structural properties of ternary solid solutions.

1. Introduction

Intriguing narrow bandgap semiconducting chalcogenide GeTe have gained extensive attention as phase-change data storage media [1, 2]. Doping is a powerful way to improve its properties [3-5]. Upon doping, the change of lattice spacing will affect its elastic constants, thermal expansion and electrical conductivity. Thus, determining the valid lattice parameter of doped GeTe is of great importance. Vegard’s law [6] has enjoyed long-lasting popularity by providing an estimation that the lattice parameters vary linearly with the solute size and concentration at a constant temperature when two constituents form a continuous solid solution, in which the solute atoms or ions distribute randomly among the solvent matrix. However, we have observed an abnormal shrinkage of GeTe when more massive W substitutes smaller Ge, which is in contrary to lattice expansion as predicted by Vegard’s law. The origin
of such counterintuitive observation remains ambiguous.

Several similar cases have been reported previously. For example, smaller Ge substitute larger Ni renders lattice dilation [7]. Fe-Ni binary alloy [8] also shows the opposite volume change, which could be explained by a numerical model based on interatomic spacings. However, this model requires that the interatomic spacings must be independent of other atoms in the neighborhood, which could be hardly fulfilled in ternary systems like W doped GeTe. In Ag-Au alloys [9], larger Ag substitute smaller Au decreases the lattice parameter due to electronic interactions between the outer electron shells of Ag and Au. This proposition provides us valuable clues that the radius of a solute atom would change by the transfer of electrons between high-energy electron shells of the solvent and the solute atoms.

Vegard’s law could be violated when factors on the electronic level show a more significant impact on the lattice parameters. In the present work, we have investigated the structure of W alloyed GeTe from the perspective of chemical bonding on the electronic scale through first-principles calculations. We attribute the anomalous shrinkage to the electronegativity effects between W and Te, which render shorter W-Te bonds and denser local structures around W.

2. Computational methods

The \textit{ab initio} calculations were performed within the framework of density functional theory (DFT) using projector augmented wave (PAW) pseudopotential method as implemented in the Vienna \textit{ab initio} simulation package (VASP) [10]. The generalized gradient approximation of Perdew-Burke-Ernzerhof (PBE) and a cutoff of 300 eV were applied for the structural relaxation and self-consistent calculations [11]. Strongly constrained and appropriately normed meta-generalized gradient approximation (SCAN) has been reported to achieve remarkable accuracy for lattice constants [12]. Thus, it was applied to confirm the anomalous volume shrinkage obtained by PBE. Gaussian smearing with a smearing width of 0.02 eV was used to determine the partial occupancies. The Gamma centered \( k \)-mesh was set to \( 4 \times 4 \times 4 \) for relaxation and \textit{ab initio} self-consistent calculations.
The initial configuration is a rock salt $2 \times 2 \times 2$ supercell of $W_{m}Ge_{32-m}Te_{32}$ ($m = 0, 1, 2, 3$) containing 64 atoms ($WGeTe$), where $W$ and $Ge$ occupy one sublattice while $Te$ ions occupy the other sublattice. $WGeTe$ models with different configurations were relaxed using the conjugate-gradient method before performing \textit{ab initio} self-consistent calculations. As full relaxations for $WGeTe$ would induce structural collapse [13], we force the models to retain their original rock salt symmetry by implementing symmetry preserved relaxations in terms of cell volume and ionic positions, respectively. The relaxations were terminated when the energy difference between iterations was less than $1.0 \times 10^{-6}$ eV. According to our calculations for $W_{m}Ge_{32-m}Te_{32}$ ($m = 2, 3$), the configurations that $W$ dopants locate near to each other would result in lower total energy than randomly distributed ones. Therefore, only energetically favorable configurations are shown in the article.

![Fig. 1 Calculated lattice parameters of W-GeTe alloys. Results calculated by PBE were shown in black squares, and SCAN results were depicted by red circles, while the lines and arrows are guides for eyes. Insets are local structures around central Ge or W for $Ge_{32}Te_{32}$, $W_{1}Ge_{31}Te_{32}$, $W_{2}Ge_{30}Te_{32}$, $W_{3}Ge_{29}Te_{32}$](image)

3. Results and Discussion

As shown in , the lattice parameter of $W$-GeTe decreases with increasing $W$ concentration calculated by both PBE and SCAN. It can be seen that the lattice parameter of pristine GeTe calculated by PBE and SCAN is 6.013 Å and 5.936 Å,
respectively. They are both in good agreement with the experimental data measured at 773 K (6.024 Å) [14]. The observation that the lattice parameter and cell volume decrease with the increasing concentration of larger radius dopants is in direct contrary to Vegard’s law. Meanwhile, the average bond length ~ 2.747 Å for W-Te bonds is shorter than Ge-Te bonds (~ 2.951 Å) in W₁Ge₃₁Te₃₂, as shown in the inset of . Hence, W would not only induce cell volume shrinkage but render compact local structures as well. Furthermore, W-Te and W-W bonds distort the rock salt structure by introducing non-orthogonal bonds. The local distortion is so severe that the rock salt lattice could hardly be preserved when W concentration reaches 9.375 at.% (W₃Ge₂₉Te₃₂) as shown in the inset of . Therefore, we abandon models with W concentration higher than W₃Ge₂₉Te₃₂ in our symmetry preserved relaxations. The short and distorted W-Te bonds would be the origin of the peculiar volume shrinkage.

Fig. 2 Charge density difference (CDD) around W in W₁Ge₃₁Te₃₂. The yellow (blue) surfaces indicate 0.0055 e/bohr³ increase (decrease) of charge density after bonds formation

As shown in , W is surrounded by blue surfaces, indicating that W loses its outer shell electrons (-0.0055 e/bohr³) after bonding with six neighboring Te. On the other hand, yellow surfaces (+0.0055 e/bohr³) are observed between W-Te, suggesting that several electrons of W transfer to Te, resulting in W-Te bonds. The not fully occupied 5p orbitals of Te play a key role in accommodating the W outer shell electrons. Moreover, the yellow surfaces around W-Te are much larger than Ge-Te, indicating
that more electrons are involved in W-Te bonding. This character results in strong W-Te interactions and hence short W-Te bonds. Therefore, we propose that W loses its outer shell electrons, which are accommodated in Te 5p orbitals when forming strong chemical interactions with Te. Consequently, the compact local structure around W and the abnormal cell volume shrinkage are observed.

Table 1 The average values of Bader charge transfer in WGeTe.

| System       | Bader charge transfer/e |
|--------------|-------------------------|
|              | W      | Ge      | Te     |
| GeTe         | ...    | 0.451   | -0.451 |
| W₁Ge₃₁Te₃₂  | 0.616  | 0.425   | -0.431 |
| W₂Ge₃₀Te₃₂  | 0.505  | 0.394   | -0.401 |
| W₃Ge₂⁹Te₃₂  | 0.505  | 0.350   | -0.364 |

Bader charge transfers can estimate electron transfer during bonds formation[15]. As shown in Table 1, W and Ge lose their electrons while Te obtains them. It is noteworthy that W loses 0.616 e in W₁Ge₃₁Te₃₂ when it is coordinated by six Te ions. On the other hand, the Bader charge transfer of W decreases when the first coordination shell includes another W (0.505 in W₂Ge₃₀Te₃₂ and W₃Ge₂⁹Te₃₂). However, W still loses more electrons than Ge, whose Bader charge transfers are among 0.350 – 0.425 e. This might be attributed to the lower electron affinity of W nucleus to its outer shell electrons, namely the 2 electrons in the 6s orbital and the 4 electrons in the 5d orbital. Bader charge transfer data have demonstrated the former proposition that W loses outer shell electrons and induces short and strong W-Te bonds, which shrink the local structure as well as decrease cell volume of WGeTe.
Further analysis of the electron localization function (ELF) unravels the bonding properties of WGeTe on the electronic scale. The values of ELF vary from 0 to 1, representing different types of chemical bonds with various strengths [16]. Both the three-dimensional display and a two-dimensional slice of ELF around W are depicted in Fig. 3. The ELF around W is below 0.68, as shown in (a), which could be reconfirmed by the blue area within the contour line indicating ELF < 0.1, as depicted in (b). Low ELF values indicate that the outer shell electrons around W are delocalized. On the other hand, the electrons are localized around Ge, forming nearly uniform covalent bonds with Te as depicted in (a, b), which is consistent with the bonding environment in pristine GeTe. Besides, as shown in (c), the highest ELF value between W-Te (∼ 0.665) lies at 0.846 Å away from Te, which is ∼ 1/3 of the W-Te bond length (2.747 Å), demonstrating that the electrons forming W-Te bonds are localized close to Te. Such observation gives another indication that Te accepts outer shell electrons of W, constructing short bonds and resulting in denser structures in WGeTe.

4. Conclusions

In summary, by ab initio calculations, we revealed anomalous volume shrinkage and compact local structure when massive W replaces smaller Ge in rock salt GeTe alloys. The electron transfer from W to Te results in shorter ionic W-Te bonds and...
thus reduced lattice parameters. We demonstrate that electronegativity difference between dopant and host atoms is crucial in determining lattice parameters of complex alloys.

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References

[1] M. Wuttig, N. Yamada, Phase-change materials for rewriteable data storage, Nature materials 6(11) (2007) 824.
[2] Z. Sun, J. Zhou, H.-K. Mao, R. Ahuja, Peierls distortion mediated reversible phase transition in GeTe under pressure, Proceedings of the National Academy of Sciences 109(16) (2012) 5948-5952.
[3] Z. Sun, S. Tian, B. Sa, Investigation of the structure and properties of rhombohedral Cu–Ge–Te alloys by ab initio calculations, Intermetallics 32 (2013) 292-296.
[4] L. Zhang, B. Sa, J. Zhou, Z. Song, Z. Sun, Atomic scale insight into the amorphous structure of Cu doped GeTe phase-change material, Journal of Applied Physics 116(15) (2014) 153501.
[5] G. Wang, J. Zhou, S.R. Elliott, Z. Sun, Role of carbon-rings in polycrystalline GeSb2Te4 phase-change material, Journal of Alloys and Compounds 782 (2019) 852-858.
[6] L. Vegard, Die konstitution der mischkristalle und die raumfüllung der atome, Zeitschrift für Physik A Hadrons and Nuclei 5(1) (1921) 17-26.
[7] H. Chessin, S. Arajs, R. Colvin, D. Miller, Paramagnetism and lattice parameters of iron-rich iron-germanium alloys, Journal of Physics and Chemistry of Solids 24(2) (1963) 261-273.
[8] H. Moreen, R. Taggart, D. Polonis, A model for the prediction of lattice
parameters of solid solutions, Metallurgical Transactions 2(1) (1971) 265-268.

[9] V. Lubarda, On the effective lattice parameter of binary alloys, Mechanics of Materials 35(1-2) (2003) 53-68.

[10] J. Hafner, Ab-initio simulations of materials using VASP: Density-functional theory and beyond, Journal of computational chemistry 29(13) (2008) 2044-2078.

[11] J.P. Perdew, Y. Wang, Accurate and simple analytic representation of the electron-gas correlation energy, Physical Review B 45(23) (1992) 13244.

[12] J. Sun, R.C. Reamsing, Y. Zhang, Z. Sun, A. Ruzsinszky, H. Peng, Z. Yang, A. Paul, U. Waghmare, X. Wu, Accurate first-principles structures and energies of diversely bonded systems from an efficient density functional, Nature chemistry 8(9) (2016) 831.

[13] D. Shin, S. Roy, T.R. Watkins, A. Shyam, Lattice mismatch modeling of aluminum alloys, Computational Materials Science 138 (2017) 149-159.

[14] H.W. Shu, S. Jaulmes, J. Flahaut, Système AsGeTe: III. Etude cristallographique d'une famille de composés à modèles structuraux communs: β-As2Te3, As4GeTe7 et As2GenTe(3+n) (n= 1à5), Journal of Solid State Chemistry 74(2) (1988) 277-286.

[15] G. Henkelman, A. Arnaldsson, H. Jónsson, A fast and robust algorithm for Bader decomposition of charge density, Computational Materials Science 36(3) (2006) 354-360.

[16] B. Silvi, A. Savin, Classification of chemical bonds based on topological analysis of electron localization functions, Nature 371(6499) (1994) 683.