Injectivity of the Cauchy-stress tensor along rank-one connected lines under strict rank-one convexity condition

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Abstract

In this note, we show that the Cauchy stress tensor $\sigma$ in nonlinear elasticity is injective along rank-one connected lines provided that the constitutive law is strictly rank-one convex. This means that $\sigma(F + \xi \otimes \eta) = \sigma(F)$ implies $\xi \otimes \eta = 0$ under strict rank-one convexity. As a consequence of this seemingly unnoticed observation, it follows that rank-one convexity and a homogeneous Cauchy stress imply that the left Cauchy-Green strain is homogeneous, as is shown in \cite{12}.

Mathematics Subject Classification: 74B20, 74G65, 26B25

Key words: rank-one convexity, nonlinear elasticity, Cauchy stress tensor, invertible stress-strain law

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1 Introduction

The search for a priori constitutive inequalities has been termed by Truesdell [20, 21] the “Hauptproblem” of nonlinear elasticity. These constitutive inequalities should guarantee reasonable physical response under all possible circumstances [17, 18.6.3]. We focus here on one of these requirements, namely rank-one convexity, and exhibit a hitherto unnoticed consequence of strict rank-one convexity in connection with the Cauchy stress tensor.

Following a definition by Ball [2, Definition 3.2], we say that $W$ is strictly rank-one convex on $GL^+(3) = \{ X \in \mathbb{R}^{3 \times 3} \mid \det X > 0 \}$ if it is strictly convex on all closed line segments in $GL^+(3)$ with end points differing by a matrix of rank one, i.e.,

$$W(F + (1 - \theta) \xi \otimes \eta) < \theta W(F) + (1 - \theta)W(F + \xi \otimes \eta) \quad (1.1)$$

for all $F \in GL^+(3)$, $\theta \in [0, 1]$ and all $\xi, \eta \in \mathbb{R}^3$ with $F + t \xi \otimes \eta \in GL^+(3)$ for all $t \in [0, 1]$, where $\xi \otimes \eta$ denotes the dyadic product. Rank-one convexity is connected to the study of wave propagation [15] or hyperbolicity of the dynamical equations of elasticity, and plays an important role in the existence and uniqueness theory for linear elastostatics and elastodynamics [14, 6, 4, 16], cf. [10]. Important criteria for the rank-one convexity of stored energy density functions were first established by Knowles and Sternberg [9], see also [11, 13, 7].

In this paper we use the Frobenius tensor norm $\|X\|^2 = \langle X, X \rangle_{\mathbb{R}^{n \times n}}$, where $\langle X, Y \rangle_{\mathbb{R}^{n \times n}}$ is the standard Euclidean scalar product on $\mathbb{R}^{n \times n}$. If no confusion can arise, we will suppress the subscripts $\mathbb{R}^{n \times n}$. The identity tensor on $\mathbb{R}^{n \times n}$ will be denoted by $\mathbb{I}$, so that $\text{tr}(X) = \langle X, \mathbb{I} \rangle$.

Rank-one convexity is preferably expressed in terms of the stored energy density $W(F)$ or as a monotonicity requirement for the first Piola-Kirchhoff stress tensor $S_1 = DW(F)$ along rank-one lines, i.e.,

$$\langle S_1(F + \xi \otimes \eta) - S_1(F), \xi \otimes \eta \rangle_{\mathbb{R}^{3 \times 3}} > 0 \quad \forall \xi \otimes \eta \neq 0, \quad \forall F \in GL^+(3), \quad (1.2)$$

which, if $W$ is twice-differentiable, turns into the well-known strong-ellipticity condition

$$D^2W(F)(\xi \otimes \eta, \xi \otimes \eta) > 0 \quad \forall \xi \otimes \eta \neq 0, \quad \forall F \in GL^+(3). \quad (1.3)$$

Since objective stored energy density functions cannot be convex in $F$ [13], the first Piola-Kirchhoff stress $S_1(F)$ will, in general, not be injective ([14, 6.2.38, 17, 18.4.5]). However, the strict monotonicity condition (1.2) means that $S_1(F + \xi \otimes \eta) = S_1(F)$ implies $\xi \otimes \eta = 0$. This motivates the following

**Definition 1.1.** The stress tensor $S$ is rank-one injective at $F$ if

$$S(F + \xi \otimes \eta) = S(F) \iff \xi \otimes \eta = 0. \quad (1.4)$$

In this sense, if the stored energy density is strictly rank-one convex, then the first Piola-Kirchhoff stress tensor $S_1(F)$ is everywhere rank-one injective.

The only well-known consequence of rank-one convexity in connection to the Cauchy stress tensor are the Baker-Ericksen inequalities [1] for the principal values of the Cauchy stress. These, however, are meaningful only for isotropy [3].

Here, we show by a short and elementary calculation that strict monotonicity of the first Piola-Kirchhoff stress tensor $S_1$ along rank-one lines implies injectivity of the Cauchy stress tensor along rank-one lines.

This elementary observation answers a question raised in a recent contribution [12]: Is it impossible for a strictly rank-one convex stored energy to admit a continuous deformation that corresponds...
to a homogeneous Cauchy stress field but has jumps in its deformation gradient field across planar interfaces? Indeed, in [12] we show that a non rank-one convex formulation may allow for a deformation with a homogeneous Cauchy stress field but an inhomogeneous left Cauchy-Green strain field.

We consider the following general situation: Let

$$\sigma: \text{GL}^+(3) \to \text{Sym}(3), \quad F \mapsto \sigma(F)$$

denote the Cauchy stress response function induced by the stored energy density $W$, and let $F \in \text{GL}^+(3)$ be such that

$$\sigma(F + \xi \otimes \eta) = \sigma(F)$$

for some $\xi \otimes \eta \neq 0$. We recall the basic relation [3]

$$\sigma(F) = S_1(F) (\text{Cof}(F))^{-1}$$

and note that in case of isotropy we may write

$$\sigma(F) = \tilde{\sigma}(F F^T) = \tilde{\sigma}(B), \quad \tilde{\sigma}: \text{Sym}^+(3) \to \text{Sym}(3), \quad B \mapsto \tilde{\sigma}(B) \, . \quad (1.7)$$

In isotropic nonlinear elasticity, a number of energies (suitable Neo-Hooke, Mooney-Rivlin [3, 14], the exponentiated Hencky energy [13]) define an invertible Cauchy stress-strain relation, in the sense that the mapping $B \mapsto \tilde{\sigma}(B)$ is invertible. In this case $\sigma(F + \xi \otimes \eta) = \tilde{\sigma}(\tilde{B}) = \tilde{\sigma}(B) = \sigma(F)$ leads to $B = \tilde{B}$. This, together with $\det \tilde{F} = \det F > 0$ implies $\xi \otimes \eta = 0$ in (1.5). A self-contained elementary proof of this fact is given in the appendix.

Our subsequent development will be independent of any invertibility assumption for the Cauchy stress $\sigma$ in the isotropic representation with $\tilde{\sigma}$.

### 2 Injectivity of the Cauchy-stress tensor along rank-one lines for strictly rank-one convex energies

We will show that equality (1.5) combined with strict rank-one convexity in the format of (1.2) leads to a contradiction.\footnote{The following alternative proof, which uses the identity $\text{Cof}(F + \xi \otimes \eta) = \text{Cof}(F \eta) $, was kindly suggested by the reviewer:}

**Proof.** To this aim, using (1.6) we compute

$$\sigma(F + \xi \otimes \eta) = S_1(F + \xi \otimes \eta) (\text{Cof}(F + \xi \otimes \eta))^{-1}, \quad \sigma(F) = S_1(F) (\text{Cof}(F))^{-1} \, . \quad (2.1)$$

Hence, from (1.5) it follows that

$$S_1(F + \xi \otimes \eta)(\text{Cof}(F + \xi \otimes \eta))^{-1} = S_1(F) (\text{Cof}(F))^{-1}$$

$$\iff$$

$$S_1(F + \xi \otimes \eta) = S_1(F) (\text{Cof}(F))^{-1} (\text{Cof}(F + \xi \otimes \eta)) \, . \quad (2.2)$$

If the stored energy density function is strictly rank-one convex, the latter identity implies that if $\sigma(F + \xi \otimes \eta) = \sigma(F)$, then $\xi \otimes \eta = 0$. 

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Using now the expansion \( \text{Cof}(AB) = \text{Cof}(A) \text{Cof}(B) \) and \( (\text{Cof}A)^{-1} = \text{Cof}(A^{-1}) \) for all \( A, B \in \text{GL}^+(3) \), we obtain
\[
S_1(F + \xi \otimes \eta) = S_1(F) \text{Cof}(F^{-1} F + F^{-1} \xi \otimes \eta) = S_1(F) \text{Cof}(\mathbb{1} + F^{-1} \xi \otimes \eta). \quad (2.3)
\]

Using now the expansion \( \text{Cof}(\mathbb{1} + H) = \text{Cof}(\mathbb{1}) + D \text{Cof}(F) \big|_{\mathbb{1}} \cdot H + \text{Cof}(H) \), see \cite{19}, we find
\[
\text{Cof}(\mathbb{1} + F^{-1} \xi \otimes \eta) = \text{Cof}(\mathbb{1}) + D \text{Cof}(F) \big|_{\mathbb{1}} \cdot (F^{-1} \xi \otimes \eta) + \underbrace{\text{Cof}(F^{-1} \xi \otimes \eta)}_{=0}, \quad (2.4)
\]
and since
\[
D \text{Cof}(F).H = \left( \langle F^{-T}, H \rangle \mathbb{1} - F^{-T} H^T \right) \text{Cof} F \Rightarrow D \text{Cof}(F) \big|_{\mathbb{1}} \cdot H = \langle \mathbb{1}, H \rangle \mathbb{1} - H^T, \quad (2.5)
\]
we can rewrite equality \( (2.2) \) as
\[
S_1(F + \xi \otimes \eta) = S_1(F) \left[ \mathbb{1} + D \text{Cof}(F) \big|_{\mathbb{1}} \cdot (F^{-1} \xi \otimes \eta) \right] = S_1(F) \left[ \mathbb{1} + (\mathbb{1}, (F^{-1} \xi \otimes \eta)) \mathbb{1} - (F^{-1} \xi \otimes \eta)^T \right]. \quad (2.6)
\]

Going back to the strict rank-one convexity condition \( (2.2) \), we compute now
\[
\langle S_1(F + \xi \otimes \eta) - S_1(F), \xi \otimes \eta \rangle = \langle S_1(F) \left[ \mathbb{1} + (\mathbb{1}, (F^{-1} \xi \otimes \eta)) \mathbb{1} - (F^{-1} \xi \otimes \eta)^T \right] - S_1(F), \xi \otimes \eta \rangle
= \langle (\mathbb{1}, (F^{-1} \xi \otimes \eta)) S_1(F) - S_1(F) (F^{-1} \xi \otimes \eta)^T, \xi \otimes \eta \rangle
= \langle (\mathbb{1}, (F^{-1} \xi \otimes \eta)) (S_1(F), \xi \otimes \eta) - (S_1(F), (\xi \otimes \eta) (F^{-1} \xi \otimes \eta)) \rangle
= \langle (\mathbb{1}, (F^{-1} \xi \otimes \eta)) (S_1(F), \xi \otimes \eta) - (S_1(F), (\eta, F^{-1} \xi) (\xi \otimes \eta)) \rangle
= \langle (\eta, F^{-1} \xi) (S_1(F), \xi \otimes \eta) - (\eta, F^{-1} \xi) (S_1(F), (\xi \otimes \eta)) \rangle = 0. \quad (2.7)
\]

Here, we have used that \( \langle \mathbb{1}, a \otimes b \rangle_{\mathbb{R}^{3 \times 3}} = \langle b, a \rangle_{\mathbb{R}^3} \) as well as \( (a \otimes b) (c \otimes d) = \langle b, c \rangle (a \otimes d) \), for all \( a, b, c, d \in \mathbb{R}^3 \).

Therefore, the assumption of the non-injectivity along rank-one lines \( (1.5) \) is in contradiction to the strict rank-one convexity \( (1.2) \).

In summary, we have shown that strict rank-one convexity implies that
\[
\sigma(F + \xi \otimes \eta) = \sigma(F) \quad \text{is impossible for a nontrivial} \quad \xi \otimes \eta \neq 0, \quad \xi, \eta \in \mathbb{R}^3. \quad (2.8)
\]

In these terms, we have thus proved that
\[
\text{strict rank-one convexity} \quad \Rightarrow \quad \text{the Cauchy stress} \quad \sigma \quad \text{is rank-one injective for all} \quad F \in \text{GL}^+(3). \quad (2.9)
\]

3 Conclusion

Our simple calculation shows that for strictly rank-one convex stored energy density functions it is impossible to have a constant Cauchy-stress field in response to a rank-one connected laminate microstructure. Our result suggests also that some form of injectivity for the Cauchy stress is natural to require in nonlinear elasticity and this injectivity should be the object of further studies.

In order to give added perspective to our result on injectivity of the Cauchy stress, let us consider the uni-constant Blatz-Ko stored energy density function
\[
W(F) = \frac{\mu}{2} \left( \|F\|^2 + \frac{2}{\det F} - 5 \right).
\]
This function is strictly polyconvex, hence strictly rank-one elliptic with Cauchy stress

\[ \tilde{\sigma} : \text{Sym}^+(3) \to \text{Sym}(3), \quad \tilde{\sigma}(B) = \frac{\mu}{\det B} \left( \sqrt{\det B \cdot B} - \mathbb{1} \right). \]

(3.10)

The Cauchy stress in \((3.10)\) is not bijective, which can be seen along the family \(B = \alpha \cdot \mathbb{1}, \alpha > 0\).

The spherical part \(\frac{1}{3} \text{tr}(\sigma)\) of the Cauchy stress first increases for increasing \(\alpha\) and then decreases.

Thus strict polyconvexity alone is not enough to prevent this unphysical response \([8]\). We need to require a condition beyond polyconvexity. Injectivity of the Cauchy stress is a candidate implying the classical pressure-compression inequality \([13]\)

\[ \frac{1}{3} \text{tr}(\sigma(\lambda \mathbb{1})) \cdot [\lambda - 1] > 0, \]

(3.11)

which would already exclude the deficiency of the Blatz-Ko strain energy.

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## 5 Appendix

In this appendix we show\(\footnote{The following alternative proof was kindly suggested by the reviewer: Rewriting (5.16) as \( (F \mathbf{\eta} + \|\mathbf{\eta}\|^2 \mathbf{\xi}) \otimes \mathbf{\xi} = -\mathbf{\xi} \otimes F \mathbf{\eta} \)}\) that

\[
\hat{F} \hat{F}^T = F F^T, \quad \hat{F} = F + \mathbf{\xi} \otimes \eta, \quad \det \hat{F}, \det F > 0 \implies \mathbf{\xi} \otimes \eta = 0. \tag{5.12}
\]

We note that \(\hat{F}\) and \(F\) are twins \[17\] Sect. 2.5] since they are rank-one connected and their principal stretches coincide. Here, not only their principal stretches coincide, but the left-stretch tensor is the same as well.

**Proof.** Since \(\hat{F} \hat{F}^T = F F^T\), we see that \((\det \hat{F})^2 = (\det F)^2\), and by assumption \(\footnote{\text{Weyl tensor}}\) we can conclude that \(\det \hat{F} = \det F\). Since

\[
\det(F + \mathbf{\xi} \otimes \eta) = \det \left( F \left( \mathbb{1} + F^{-1} \mathbf{\xi} \otimes \eta \right) \right) = \det F \cdot \det(\mathbb{1} + F^{-1} \mathbf{\xi} \otimes \eta) = \det F \cdot \left( 1 + \text{tr}(F^{-1} \mathbf{\xi} \otimes \eta) \right) \tag{5.13}
\]

and \(\det(F + \mathbf{\xi} \otimes \eta) = \det \hat{F} = \det F\), by \(\footnote{\text{Weyl tensor}}\) we conclude from \(\footnote{\text{Weyl tensor}}\)

\[
\text{tr}(F^{-1} \mathbf{\xi} \otimes \eta) = \langle F^{-1} \mathbf{\xi}, \mathbf{\eta} \rangle = 0. \tag{5.14}
\]

Assumption \(\footnote{\text{Weyl tensor}}\) and \(\footnote{\text{Weyl tensor}}\) together imply

\[
\hat{F} \hat{F}^T = F F^T + F \mathbf{\eta} \otimes \mathbf{\xi} + \mathbf{\xi} \otimes F \mathbf{\eta} + \|\mathbf{\eta}\|^2 (\mathbf{\xi} \otimes \mathbf{\xi}) = F F^T, \tag{5.15}
\]

thus we must have

\[
F \mathbf{\eta} \otimes \mathbf{\xi} + \mathbf{\xi} \otimes F \mathbf{\eta} + \|\mathbf{\eta}\|^2 (\mathbf{\xi} \otimes \mathbf{\xi}) = 0. \tag{5.16}
\]

We introduce \(\hat{\mathbf{\xi}} = F^{-1} \mathbf{\xi}, \mathbf{\xi} = F \hat{\mathbf{\xi}}\) and insert into \(\footnote{\text{Weyl tensor}}\) and \(\footnote{\text{Weyl tensor}}\) to yield

\[
F \mathbf{\eta} \otimes F \hat{\mathbf{\xi}} + F \hat{\mathbf{\xi}} \otimes F \mathbf{\eta} + \|\mathbf{\eta}\|^2 (F \hat{\mathbf{\xi}} \otimes F \hat{\mathbf{\xi}}) = 0, \quad \langle \hat{\mathbf{\xi}}, \mathbf{\eta} \rangle = 0. \tag{5.17}
\]

This is equivalent to

\[
F \left\{ \mathbf{\eta} \otimes \hat{\mathbf{\xi}} + \hat{\mathbf{\xi}} \otimes \mathbf{\eta} + \|\mathbf{\eta}\|^2 (\hat{\mathbf{\xi}} \otimes \hat{\mathbf{\xi}}) \right\} F^T = 0, \quad \langle \hat{\mathbf{\xi}}, \mathbf{\eta} \rangle = 0. \tag{5.18}
\]

Since \(\det F > 0\) we have as well

\[
\mathbf{\eta} \otimes \hat{\mathbf{\xi}} + \hat{\mathbf{\xi}} \otimes \mathbf{\eta} + \|\mathbf{\eta}\|^2 (\hat{\mathbf{\xi}} \otimes \hat{\mathbf{\xi}}) = 0, \quad \langle \hat{\mathbf{\xi}}, \mathbf{\eta} \rangle = 0. \tag{5.19}
\]

Multiplying \(\footnote{\text{Weyl tensor}}\) with \(\mathbf{\eta} \neq 0\) we obtain \(\mathbf{\eta} \langle \hat{\mathbf{\xi}}, \mathbf{\eta} \rangle + \hat{\mathbf{\xi}} \|\mathbf{\eta}\|^2 + \|\mathbf{\eta}\|^2 \hat{\mathbf{\xi}} \langle \hat{\mathbf{\xi}}, \mathbf{\eta} \rangle = 0\). Hence, \(\|\mathbf{\eta}\|^2 = 0\) implies \(\hat{\mathbf{\xi}} = 0\). \(\blacksquare\)