A Hot Ultraviolet Flare on the M Dwarf Star GJ 674

Cynthia S. Froning1, Adam Kowalski2,3,4, Kevin France2,3, R. O. Parke Loyd5, P. Christian Schneider6, Allison Youngblood7,12, David Wilson1, Alexander Brown8, Zachary Berta-Thompson8, J. Sebastian Pineda8, Jeffrey Linsky9, Sarah Rugheimer10, and Yamila Miguel11

1 McDonald Observatory, University of Texas at Austin, Austin, TX 78712, USA; cfroning@astro.as.utexas.edu
2 Laboratory for Astrophysical and Space Astronomy, University of Colorado, 600 UCB, Boulder, CO 80309, USA
3 Department of Astrophysical and Space Astronomy, University of Colorado, 389 UCB, Boulder, CO 80309, USA
4 National Solar Observatory, University of Colorado at Boulder, 3665 Discovery Drive, Boulder, CO 80303, USA
5 School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA
6 Hamburger Sternwarte, Gojenbergsweg 112, D-21029 Hamburg, Germany
7 Goddard Space Flight Center, Greenbelt, MD 20771, USA
8 Center for Astrophysics and Space Astronomy, 389 UCB, University of Colorado, Boulder, CO 80309, USA
9 JILA, 440 UCB, University of Colorado, Boulder, CO 80309
10 University of Oxford, Clarendon Laboratory, AOPP, Sherrington Road, Oxford OX1 3PU, UK
11 Leiden Observatory, P.O. Box 9500, 2300 RA Leiden, The Netherlands

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Abstract

As part of the Mega-Measurements of the Ultraviolet Spectral Characteristics of Low-Mass Exoplanetary Systems Hubble Space Telescope (HST) Treasury program, we obtained time-series ultraviolet spectroscopy of the M2.5V star, GJ 674. During the far-ultraviolet (FUV) monitoring observations, the target exhibited several small flares and one large flare ($E_{\text{FUV}} = 10^{30.75}$ erg) that persisted over the entirety of an HST orbit and had an equivalent duration >30,000 s, comparable to the highest relative amplitude event previously recorded in the FUV. The flare spectrum exhibited enhanced line emission from chromospheric, transition region, and coronal transitions and a blue FUV continuum with an unprecedented color temperature of $T_C \approx 40,000 \pm 10,000$ K. In this Letter, we compare the flare FUV continuum emission with parameterizations of radiative hydrodynamic model atmospheres of M stars. We find that the observed flare continuum can be reproduced using flare models but only with the ad hoc addition of a hot, dense emitting component. This observation demonstrates that flares with hot FUV continuum temperatures and significant extreme-ultraviolet/FUV energy deposition will continue to be of importance to exoplanet atmospheric chemistry and heating, even as the host M dwarfs age beyond their most active evolutionary phases.

Key words: planets and satellites: atmospheres – stars: activity – stars: chromospheres – stars: flare – stars: individual (GJ 674) – stars: low-mass

1. Introduction

Active low-mass M (dMe) stars are subject to frequent, energetic stellar flares. The energies can exceed values seen in solar flares by factor of up to $1 \times 10^6$ on the most active M dwarfs (Osten et al. 2010, 2016), and large flares are observed on relatively inactive M dwarfs on occasion as well (Paulson et al. 2006). During the impulsive phase of the flare, the stellar atmosphere is believed to be heated by a beam of mildly relativistic, nonthermal electrons accelerated after magnetic field reconnection and directed from the corona along the field lines into the stellar chromosphere. The electron beam drives a downflowing compression of the chromosphere; these “chromospheric condensations” (CCs) are possible sources of the visible band signatures of M-star flares. Radiative-hydrodynamic (RHD) simulations of the CCs can match the observed characteristics of dMe flares by enhancing the levels of electron beam heating over solar models, but such models require the introduction of multiple emitting regions on the star, including regions with lower densities, to provide reasonable Balmer line broadening consistent with observations. They also require flux density levels that may drive beam instabilities, suggesting that alternate heating mechanisms may be at play (Kowalski et al. 2015, 2017).

Scientists investigating flare physics in dMe stars have long pursued dedicated observing programs, but the field is also benefiting from the increased interest in the spectral properties of exoplanet host stars, particularly from the expanding observational sample of “optically inactive” M-dwarf stars. One such survey, the Mega-Measurements of the Ultraviolet Spectral Characteristics of Low-Mass Exoplanetary Systems (Mega-MUSCLES) Treasury Survey, is probing the high energy radiation environment of M stars. Mega-MUSCLES and its predecessor, MUSCLES, use panchromatic (5 Å–5.5 μm) spectroscopy and time-series monitoring of M stars to characterize the stellar energetic radiation environment and activity levels, and their effects on exoplanet atmospheres and habitability (France et al. 2013, 2016; Loyd et al. 2016, 2018; Youngblood et al. 2016, 2017).

As part of the Mega-MUSCLES survey, we observed the M2.5V ($0.35 M_\odot$) star, GJ 674. It hosts a $m \sin i = 11.09 M_\oplus$ planet in a 4.69 day (0.039 au) orbit (Bonfils et al. 2007). GJ 674 is classified as a weakly active star: while there is no emission signature in Hα characteristic of flare stars, there is an emission reversal in Ca II and signs of regular spot activity. The MUSCLES survey has shown that even optically inactive M stars exhibit frequent ultraviolet flare activity (France et al. 2016; Loyd et al. 2018). GJ 674 is no exception. During five monitoring orbits, Hubble Space Telescope (HST)
far-ultraviolet (FUV) observations caught several small flares and one large one. The large flare exhibited strong emission in lines tracing chromospheric and coronal regions of the upper atmosphere of the star and blue FUV continuum emission with a hot ($T_c \approx 40,000$ K) color temperature.

The latter phenomenon is the motivation for this work. In this Letter, we present the properties of this unique flare observation and investigate models to reproduce them. In Section 2, we describe the observations and data analysis. Section 3 presents the RHD model parameterizations and fits to the flare FUV flux. Section 4 discusses the implications of the results for our understanding of flare physics and the effects of such flares on exoplanet atmospheres in M stars.

2. Observations

We observed GJ 674 in 2018 April 2–3 using the Cosmic Origins Spectrograph (COS) on HST (Green et al. 2012). The visit included five HST orbits in COS G130M, providing time-resolved FUV spectroscopy of the star. The COS G130M/1222 mode has spectral coverage from 1065–1365 Å at resolving powers of $R > 10,000$. The time-series spectral light curves were extracted from the time-tag event files for each exposure by summing the counts around the location of the spectral trace in the 2D event list and subtracting background regions of the same size offset in the cross-dispersion direction.

The COS light curve and the quiescent and flare spectra are shown in Figure 1. Several small flares were observed during the visit, as well as one large flare in the last orbit. We defined the flare spectrum as the entirety of orbit five and the out-of-flare spectrum as the average of the previous four orbits. The FUV flare energy was calculated as per Equation (1) in Loyd et al. (2018) and was found to be $E_{FUV} = 10^{30.75}$ erg. If defined in terms of the equivalent duration, $\delta$, that the star would have to emit in quiescence to match the observed flare energies, the flare had $\delta \approx 30,000$ s. The equivalent duration of the COS flare is $2.0–4.4 \times$, which are the largest values seen in previous HST FUV observations of M-star flares (Loyd et al. 2018, 2019). More energetic flares (in absolute emission energy) have been observed, but because GJ 674 is not a “flare star” (according to its lack of Hα emission in quiescence; Kowalski et al. 2009), the COS flare is substantially more energetic relative to the quiescent flux of the star than previously observed flares. In terms of equivalent duration, the flare rivals the Great Flare observed on AD Leo by the International Ultraviolet Explorer (IUE; Hawley & Pettersen 1991; Loyd et al. 2018, $\delta \approx 40,000$ s).

Figure 2 highlights several interesting spectral regions in and out of the flare. During the flare, line emission was enhanced in a number of lines tracing chromospheric and transition region emission (e.g., C II, C III, Si III, Si IV, N V; with formation temperatures ranging from $T_{\text{form}} \approx 3 \times 10^4$ K to $T_{\text{form}} \approx 2 \times 10^5$ K) as well as lines tracing hotter coronal regions (Fe XII, Fe XIX, Fe XXI; $T_{\text{form}} \approx 10^5–10^6$ K). In general, the lower temperature chromospheric lines showed similar peak-to-quiet-flux ratios to the largest flare in the MUSCLES sample of optically inactive M stars (on GJ 876; France et al. 2016), with C III and Si III increasing in flux by a factor of $\approx 100$ and C II by 30–50. However, the higher formation temperature N V line and the FUV continuum emission increased by larger factors in the GJ 674 flare—$\approx 10$ versus $\approx 5$ for the GJ 876 flare—and the [Fe XXI] λ1354 line, only marginally detected in the GJ 876 flare, increased by a factor of $\approx 20$ here.

In Table 1, we summarize the emission properties of the flare, giving the fluxes, line widths, and velocities for a selection of transitions before the flare and integrated over the flare. We also give the FUV continuum flux at the flare peak, the integrated flare energy in the FUV, and estimates of the integrated extreme-ultraviolet (EUV) flare energy. For the latter, we directly scaled the FUV emission to the EUV using the semi-empirical M-dwarf stellar model from Fontenla et al. (2016) and also calculated the EUV energy using the relation between the strength of the N V line and the EUV flux from France et al. (2018). (Note that both of these estimates rely on quiescent scaling relations applied to the flare spectra.)

The emission line flare properties are also presented as inputs for future modeling. Here, we note a few features displayed during the flare. The N V doublet line ratio indicates that the line is optically thin both in quiescence and during the flare, reflecting its origin in the transition region. The hot coronal transition [Fe XXI] exhibits a blueshift of $\approx -25$ km s$^{-1}$ (relative to its quiescent velocity) during the flare, indicative of hot coronal plasma being pushed outward by the exploding plasma below. In the chromospheric lines, the broad components of the lines widen by $\approx 5–25$ km s$^{-1}$, which may indicate an increase in the turbulent velocity of the line-forming region during the flare. The broad line components also redshift by $\approx 5$ km s$^{-1}$, while the narrow components remain stationary.
We tracked the time evolution of the emission lines and the continuum during the flare, summarized in the left panel of Figure 3. The light curves have been normalized to the peak flux in each line. Notably, the FUV continuum traces the time evolution of the chromospheric lines, peaking earlier and declining faster than the coronal [Fe XXI] emission. The latter remains high well after the impulsive phase of the flare, indicating ongoing heating in the corona. We fit a single temperature blackbody to the FUV flare continuum emission. The best fit, shown in Figure 3, has a color temperature of \( T \approx 40,000 \pm 10,000 \) K. It is possible that the FUV continuum contains contributions from unresolved recombination continua and a superposition of discrete emission lines. However, the HST-COS data analyzed here are of sufficient depth and spectral resolution to separate and resolve neutral and singly ionized atomic species. Additionally, bins that showed evidence of discrete spectral features or rms flux uncertainty >50% of the bin flux were discarded.

### 3. Model Fits

RHD models of chromospheres predict the formation of a hot, dense CC layer caused by electron beam heating. When the beam energy flux rate is high, the flaring regions exceed a continuum optical depth of \( \tau \approx 1 \) and produce a spectrum with a characteristic temperature \( T \approx 10^4 \) K with a small Balmer jump, consistent with the signatures of flares seen in optical observations of dMe stars (e.g., Hawley & Fisher 1992; Fuhrmeister et al. 2011; Kowalski et al. 2013, 2016). In a recent paper, Kowalski & Allred (2018, hereafter K18) developed parameterizations of CC properties. The parameterizations can be used to predict observables from existing RHD model flare atmospheres without generating computationally intensive full simulations. Using a suite of existing RHD models, K18 developed prescriptions for the continuum optical depth and emergent flux as a function of \( T_{\text{ref}} \) and \( m_{\text{ref}} \), the temperature and column mass (\( m \) is in units of g cm\(^{-2}\)) at the chromosphere depth in which the chromospheric gas is in steady state with the electron beam heating.

We use these parameterizations to investigate modifications to the RHD models needed to fit the hot, blue FUV continuum emission seen in the GJ 674 flare. RHD models with very high electron beam energy fluxes of \( 10^{13} \) erg cm\(^{-2}\) s\(^{-1}\) can generate the \( T \approx 10^4 \) K thermal component seen in previously observed flares but not the blue FUV continuum observed here, as the models turn over in the near ultraviolet (NUV; Kowalski et al. 2015). Instead, Figure 4 shows an example composite model selected to be consistent with the observed FUV continuum during the COS flare. The first model (dotted line) is a very hot, dense \( (T_{\text{ref}} = 75,000 \) K, log \( m_{\text{ref}} = -1.7 \) emitting slab of material, generated from a CC model consisting only of the downflowing condensation layer (emission from the stationary layer has been set to zero). The
Table 1
Line and Continuum Fluxes for the COS Flare

| Line       | Quiescent Flux \(10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}\) | Velocity \(\text{km s}^{-1}\) | FWHM \(\text{km s}^{-1}\) | Flare Peak Flux \(10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}\) | Velocity \(\text{km s}^{-1}\) | FWHM \(\text{km s}^{-1}\) |
|------------|--------------------------------------------------|----------------------------|-----------------|--------------------------------------------------|----------------------------|-----------------|
| Si iii 1206 n | 6.0 ± 0.4                                         | 2.3 ± 0.5                  | 33 ± 2          | 66 ± 1                                           | 3.2 ± 0.2                  | 37 ± 1          |
| Si iii 1206 b | 4.3 ± 0.6                                         | 7 ± 1                      | 101 ± 5         | 44 ± 2                                           | 12.4 ± 0.5                 | 110 ± 1         |
| N V 1238 n   | 6.0 ± 0.4                                         | −10.6 ± 0.3                | 33 ± 2          | 21 ± 2                                           | −11.6 ± 0.3                | 29 ± 1          |
| N V 1238 b   | 3.0 ± 0.9                                         | −11.4 ± 0.7                | 64 ± 5          | 17 ± 2                                           | −7.2 ± 0.6                 | 70 ± 2          |
| N V 1242 n   | 2.6 ± 0.3                                         | −4.5 ± 0.4                 | 32 ± 2          | 14 ± 6                                           | −9.1 ± 0.3                 | 38 ± 1          |
| N V 1242 b   | 1.9 ± 0.6                                         | −5 ± 1                     | 64 ± 6          | 4.1 ± 0.8                                        | 3 ± 3                      | 97 ± 7          |
| C II 1334 n  | 5.3 ± 0.2                                         | −16.4 ± 0.2                | 25 ± 1          | 32.9 ± 0.9                                       | −17.2 ± 0.2                | 28.7 ± 0.6      |
| C II 1334 b  | 2.6 ± 0.3                                         | −11.8 ± 0.8                | 75 ± 3          | 23 ± 1                                           | −7.3 ± 0.6                 | 85 ± 2          |
| C II 1335 n  | 9.1 ± 0.4                                         | −15.3 ± 0.2                | 30.8 ± 0.9      | 61 ± 1                                           | −17.3 ± 0.2                | 41.7 ± 0.6      |
| C III 2965 b | 4.6 ± 0.6                                         | 13.0 ± 0.6                 | 75 ± 3          | 28 ± 2                                           | −9.8 ± 0.6                 | 98 ± 2          |
| [Fe XII] 1242| 0.48 ± 0.02                                       | −14.9 ± 0.6                | 50 ± 2          | ...                                             | ...                       | ...             |
| [Fe XIX] 1118| 0.31 ± 0.05                                       | −21 ± 4                    | 106 ± 12        | 2.6 ± 0.3                                        | −17 ± 3                    | 100 ± 8         |
| [Fe XXI] 1354| 0.21 ± 0.06                                       | 12 ± 8                     | 101 ± 26        | 4.7 ± 0.2                                        | −13 ± 2                    | 125 ± 4         |

Notes. All emission lines, with the exception of the Fe lines (for which the signal-to-noise ratio (S/N) was too low) are fit with double Gaussian profiles with broad (b) and narrow (n) components, consistent with typical fits to M-dwarf FUV emission lines (Wood et al. 1997). The Fe lines are fit with single Gaussians.

a The [Fe XII] line cannot be fit reliably in the flare spectrum because it is blended with the broad wings of Lyα. We can place a conservative estimate on the increase in the peak (not integrated) line emission during the flare of <50% relative to its quiescent flux.

b The FUV continuum flux density at the peak of the flare.

c The integrated flare energy in the FUV, excluding detector gaps and airglow lines, summed over the entirety of orbit five.

d Estimated by scaling the semi-empirical M-dwarf spectrum of Fontenla et al. (2016) to the measured FUV flare energy.

e Estimated using the relations between N V emission and EUV flux from France et al. (2018). The calculation of the stellar bolometric flux gave \(F_{\text{bol}} = 2.3 \times 10^{-6} \text{ erg cm}^{-2} \text{ s}^{-1}\), assuming \(d = 4.5 \text{ pc}\), \(R_s = 0.35 R_\odot\), and \(T_{\text{eff}} = 3400 \text{ K}\) (Pecaut & Mamajek 2013; Newton et al. 2016). Then, \(F_{\nu,v,\text{flare}}/F_{\text{bol}} = 2.5 \times 10^{-6}\) and \(F(90–911 \text{ Å, flare}) = 9.6 \times 10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}\). The overall EUV flare energy, \(E(90–911 \text{ Å})\) relies on the Fontenla et al. (2016) models to scale the Lyman continuum EUV component to the total EUV emission.

Figure 3. The left panel displays the evolution of the COS flare in several emission lines. The light curves have been normalized to 1.0 at their peak fluxes. The right panel shows a single temperature blackbody fit to the COS flare spectrum of GJ 674. The method of identifying the continuum and fitting follows that of France et al. (2014); continuum-only regions were verified against those identified in Loyd et al. (2018) from high S/N flare spectra. The points in orange were used in the fit, while the broad Lyα emission line wings and the short wavelength end of the waveband where the S/N was low (shown in gray) were excluded. The errors on the data points are the 1σ uncertainties in the background-subtracted FUV-continuum counts in 1 Å bins.

second model (dashed line) is a CC parameterization that employs a value of \(m_{\text{eff}} = -3.0\), characteristic of a F11 (10^{11} \text{ erg cm}^{-2} \text{ s}^{-1}) electron beam. It has been scaled up by a factor of 200 such that, when superposed with the hotter component, the combination spectrum exhibits a color temperature in the optical (\(\lambda > 4000 \text{ Å}\)) of 9000 K and a moderate Balmer jump ratio (Kowalski et al. 2018). Adding the two models (blue curve) results in a FUV color temperature of \(T_C \approx 26,000 \text{ K}\). If the F11 model is not included, the hot component alone produces optical temperatures and Balmer
jump ratios ($T_c \approx 12,000$ K and $\chi_{\text{flare}} = 1.7$) that are only seen in the most extreme dMe flares and may be unlikely for the less active GJ 674. Thus, while we can match the observed FUV flare continuum, it requires the addition of a hot, dense CC component not predicted by current RHD model simulations.

### 4. Discussion

With a rotation period of 33.4 days and a likely age of 0.1 to a few Gyr, GJ 674 sits at an intermediate position between rapidly rotating, young dMe stars ($P_{\text{rot}} < 10$ days, ages $< 2$ Gyr) and older, spun-down stars ($P_{\text{rot}} > 70$ days, ages of several Gyr Bonfils et al. 2007; Kiraga & Stepień 2007; Newton et al. 2016). The lack of Hα in emission and a $P_{\text{rot}} > 30$ days put GJ 674 on the inactive side of the mass-period relation for a M2.5V, 0.35 $M_\odot$ star (Newton et al. 2018). In the FUV, however, GJ 674 was active, a result that is consistent with the analysis of Loyd et al. (2018, 2019), which demonstrated that, when normalized by the basal flux of the star, dMe and inactive M stars show the same FUV flare energy-frequency relation. The large flare that we observed is unprecedented in several respects. In terms of the absolute flare energy, it is not particularly strong compared to historic dMe flare observations. However, the equivalent duration of the flare places it in the realm of the highest amplitude flares observed. Moreover, the color temperature of the continuum, fit with a blackbody of $T_c \approx 40,000$ K, is this hottest spectroscopic color temperature contribution to a broadband flare spectrum measured to date.

While direct constraints from FUV observations of M-dwarf flares have been relatively scarce to date, there have been previous indications of the presence of very hot, dense gas emission during flares. The recent large flare captured by the HAbitable Zones and M dwarf Activity across Time (HAZMAT) team had an FUV (1170–1430 Å) continuum fit by a 15,500 K blackbody; while the continuum emission in that flare was red in the FUV, its color temperature was still higher than those usually seen in optical flare spectra (Kowalski et al. 2013; Loyd et al. 2019). FUV spectra taken during the AD Leo Great Flare also showed a flat UV continuum (Hawley & Pettersen 1991). Photometric observations of flares in the FUV (1344–1786 Å) and NUV (1771–2831 Å) channels with Galaxy Evolution Explorer (GALEX) have also suggested the presence of two distinct flare components, one hotter and of shorter duration, that reaches $T_c = 50,000$ K at its peak (Robinson et al. 2005; Welsh et al. 2006).

Regarding the origin of the FUV continuum, its properties indicate that a higher column mass is heated than predicted by any RHD beam model so far explored. Our investigation shows that a combination of a CC and an additional hot chromospheric emitter could replicate the FUV continuum temperature while providing reasonable values for the optical color temperature and Balmer jump ratio.

A coronal origin, from thermal 30–50 MK gas, may also be possible but generating the requisite gas density is challenging. For example, Caspi & Lin (2010) also found two components of X-ray continuum emission in an energetic (X4.8) solar flare: a super-hot ($T_e > 30$ MK) component, likely due to direct coronal heating and compression during the magnetic reconnection process (Longcope & Guidoni 2011; Longcope et al. 2016), and a hot ($T_e \lesssim 25$ MK) component that they attributed to chromospheric evaporation. However, an extension of their model spectra to the FUV predicts $F_{142} = 5 \times 10^{-17}$ erg cm$^{-2}$ s$^{-1}$ Å$^{-1}$, a factor of $\sim 20$ below the observed FUV flux. The FUV continuum emission also followed the rapid rise and fast initial decay of the

![Flare model flux densities in the FUV and optical.](image)

**Figure 4.** Flare model flux densities in the FUV and optical. The dashed line is parameterization of a F11 solar flare model of K18 multiplied by a factor of 200 to match the optical color temperature and Balmer jump parameters seen in dMe flares. The dotted line is a $T_{\text{eff}} = 75,000$ K, log $m_{\text{ref}} = -1.7$ downflowing compression layer (the stationary flare layers are set to zero). The solid blue line is the superposition of the previous two models. The $\chi_{\text{flare}}$ values give the size of the Balmer jump ratio ($= F_{1350}/F_{1410}$).

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chromospheric lines more closely than the delayed rise and long persistence of the coronal emission. Further observations of optically inactive M stars in the optical could be used to constrain these models based on optical Balmer jump and color temperature measurements, as well as constraining the presence or absence of multiple emitting components.

With respect to exoplanets around inactive M stars, the GJ 674 large flare illustrates the continued impact of stellar irradiance on planetary atmospheres well after the runaway greenhouse phase on the pre-main sequence and the young, active dMe stage of the star. Ultraviolet (UV) continuum fluxes with color temperatures in excess of 40,000 K during flares can increase the photolysis rate by 1–2 orders of magnitude over a canonical 9000 K flare continuum temperature (Lloyd et al. 2018). The EUV/FUV energy depositions onto the surface of the planet during the flare has a particular effect on molecules such as H₂, N₂, CO, and CO₂ whose photodissociation cross sections peak <1200 Å, as well as H₂O in small gas planets (Miguel et al. 2015; Loyd et al. 2016; Venot et al. 2016). Thus, the impacts of the high UV flux emitted during flares on the chemistry and potential habitability of M-star exoplanets will continue to be of importance even as the stars age beyond their most active evolutionary phases.

5. Conclusions

We observed the M2.5V exoplanet host star, GJ 674, using COS on HST. During the FUV spectroscopic observations, the star had several small flares and one large one. The large flare spanned more than a HST orbit, had a FUV flare energy \( E_{\text{FUV}} \geq 10^{37} \) erg, and a flare equivalent duration of \( \delta \geq 30,000 \) s, making the energy of the flare relative to the quiescent stellar flux comparable to the most energetic flare previously observed in the FUV. The flare spectrum displayed enhanced emission in lines tracing chromospheric and coronal temperature regimes and a strong, blue continuum with a color temperature of \( \approx 40,000 \) K, the hottest spectroscopic color temperature contribution to broadband observations of M-star flares to date. We investigated modifications to existing RHD flare models needed to match the observed FUV continuum. We find that the addition of a very hot compression slab in the CC region can generate the requisite FUV color temperatures, but it requires gas column masses that are significantly higher than are produced in current RHD models of stellar flares. Finally, we note the effects of flares with strong, hot FUV/ EUV continuum emission on the atmospheres of exoplanets around M stars, particularly the fact that stars that have evolved past the pre-main sequence and young, active stellar phases can still drive significant atmospheric heating and photolysis in exoplanet atmospheres via continued energetic flares.

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Facility: HST (COS).

Software: Astropy (The Astropy Collaboration et al. 2018), Jupyter (Klyuyver et al. 2016), Matplotlib (Hunter 2007), NumPy (van der Walt et al. 2011).

ORCID iDs

Cynthia S. Froning https://orcid.org/0000-0001-8499-2892
Adam Kowalski https://orcid.org/0000-0001-7458-1176
Kevin France https://orcid.org/0000-0002-1002-3674
R. O. Parke Loyd https://orcid.org/0000-0001-5646-6668
Alexander Brown https://orcid.org/0000-0003-2631-3905
J. Sebastian Pineda https://orcid.org/0000-0002-4489-0135
Jeffrey Linsky https://orcid.org/0000-0003-4446-3181

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