Understanding the collective behaviour of electrons in solids is increasingly desirable as these electronic interactions give rise to many intriguing phenomena in condensed matter physics, such as superconductivity and quantum Hall effects. The nature of electrons in solids is described primarily by three quantum parameters: energy ($E$), momentum ($k$) and spin ($S$). Owing to its unique capability to probe the energy and momentum of electrons directly and simultaneously, angle-resolved photoemission spectroscopy (ARPES) has a leading role in achieving a comprehensive understanding of the electronic properties of solid-state materials.

With the rapid development of electron spectrometers, modern synchrotrons and laser light sources over the past three decades, ARPES has experienced a renaissance. The notable improvement in energy and momentum resolution with laser light sources not only enables fine features of electronic states, such as the energy gap in superconductors, to be measured with unprecedented precision, but also makes it possible to obtain information on many-body interactions in strongly correlated systems. The use of continuously tunable soft X-rays greatly enhances the bulk sensitivity of ARPES, which will have a central role in probing the electronic structure of microscale and nanoscale materials as well as materials with phase separation or multiple domains. Finally, the implementation of time-resolved ARPES with ultrafast lasers or X-ray sources makes it possible to study ultrafast electronic dynamics in the time domain and enables the unoccupied states above the Fermi level to be probed. The advances in ARPES have made it an effective and ideal tool for the identification of the unique bulk and surface electronic structures of topological materials. Topological materials thus serve as excellent examples to illustrate the capability of different ARPES techniques.

Three-dimensional topological materials are usually characterized by the nature of the surface states induced by the topology of the bulk band structure, which can be divided into two groups with respect to the bulk bandgap: insulators with non-zero bandgaps (for example, topological insulators and topological crystalline insulators) and semimetals with no bandgaps (for example, Dirac semimetals and Weyl semimetals). Taking advantage of different ARPES techniques, such as vacuum ultraviolet (VUV, with...
Key points

- Topological materials are characterized by non-trivial bulk and surface electronic states, which can be detected and distinguished by angle-resolved photoemission spectroscopy (ARPES).
- Synchrotron-based vacuum ultraviolet and soft-X-ray light make it possible to distinguish surface and bulk states through photon-energy-dependent ARPES measurements.
- The integration of spin detectors into ARPES photoelectron spectrometers enables the detection and quantification of spin polarization in band structures.
- Time-resolved ARPES with femtosecond laser pulses facilitates the study of ultrafast electronic dynamics and states above the chemical potential.
- Spatially resolved ARPES with sub-micrometre spatial resolution can be used to probe the electronic structure of microscale and nanoscale materials as well as materials with phase separation or multiple domains.

Fundamentals of ARPES

The basics of ARPES have been described in detail elsewhere\(^\text{[1-10]}\). In the following section, we provide a brief introduction to salient points that are useful in understanding ARPES studies of topological materials.

ARPES is based on the photoelectric effect, which was first discovered by Hertz in 1887 (REF\(^\text{[106]}\)); the microscopic mechanism of this effect was later explained by Einstein in 1905 (REF\(^\text{[107]}\)), when he introduced the concept of a quantum of light — the photon. In a typical ARPES measurement, a sample is placed under ultrahigh vacuum next to an electron analyser (Fig. 1a). When light is incident on the sample, electrons in the material absorb photons, and if the energy of the absorbed photons is greater than the work function of the material, the electrons can escape into the vacuum. These photoemitted electrons, known as photoelectrons, are then collected and analysed with respect to their kinetic energy and emission angle by a spectrometer. The energy and momentum of the electrons inside the sample are directly connected to those of the photoelectrons by the conservation of energy and momentum parallel to the sample surface. Under the emission angles \(\theta\) and \(\varphi\) defined in Fig. 1a, the following relationships hold (Eq. 1 and 2).

\[
E_{\text{kin}} = h\nu - \Phi - E_B \quad (1)
\]

Here, \(E_{\text{kin}}\) is the kinetic energy of a photoelectron; \(h\nu\) is the photon energy; \(E_B\) is the binding energy of the electron inside the sample; and \(\Phi\) is the work function of the material, which is the energy required for an electron at \(E_B\) to escape to the vacuum level (\(E_{\text{vac}}\)) (where \(\Phi = E_{\text{vac}} - E_B\); Fig. 1b).

Under the framework of the three-step model and the sudden approximation, the ARPES photoemission intensity \(I(k, E)\) can be written as

\[
I(k, E) = I_0(k, \nu, A)A(k, E)f(E, T) \quad (4)
\]

where \(A(k, E)\) is the one-particle spectral function from which it is possible to extract information about the
quasiparticle self-energy that encodes the band structure and correlation effects, and \( \langle E, T \rangle \) is the Fermi–Dirac distribution (where \( T \) is the temperature). The first term on the right-hand side \( f_0(k, v, A) \propto \Sigma_k \left| M_{k f} \right| \left\langle \phi_{f} \right| A \cdot P \phi_{i} \rangle \) is the photoemission matrix element, which describes the transition of the initial state \( \phi_{i} \) to the final state \( \phi_{f} \). \( P \) is the electronic momentum operator, and \( A \) is the electromagnetic vector potential, which depends on the photon polarization and energy. The matrix element carries no direct information on the band dispersion. However, the matrix element can provide important orbital information on electronic states if specific measurement geometries are implemented (see Supplementary Information for further information) and has been used widely in the study of, for example, iron-based superconductors\(^{107-109}\).

**VUV ARPES**

An important parameter of ARPES is the incident photon energy. In the past few years, the energy range of photons has greatly expanded, owing to the development of laser and synchrotron light sources. At present, the incident light used varies from VUV light to soft and even hard X-rays. Among these, the most commonly used is VUV light.

The universal curve of the inelastic mean free path of a photoelectron as a function of its kinetic energy\(^{10}\) (Fig. 2a) shows that for incident photon energies above 20 eV in the VUV region, the mean free path is short (<0.6 nm).

Thus, ARPES is an extremely surface-sensitive technique, and a considerable fraction of the total ARPES signal will be representative of only the topmost surface layer of a sample. This surface sensitivity is an advantage when the aim is to probe surface states, such as those in topological insulators and Weyl semimetals, but requires ARPES experiments to be performed on atomically clean and well-ordered flat surfaces. To obtain a clean surface and to avoid surface contamination, single crystals are typically cleaved in situ and measured in ultrahigh vacuum chambers, but even then, the freshly cleaved surface has a finite measurement lifetime. To maximize the sample measurement time, it is necessary to use a vacuum with pressures lower than \( 5 \times 10^{-11} \) Torr, which can be routinely achieved by using modern ultrahigh-vacuum systems. Although flat mirror-like surfaces can be obtained through straightforward cleaving for many materials, there are numerous materials that are not cleavable (especially 3D materials), which limits the applicability of ARPES. As an alternative to cleaving, freshly prepared thin films can be used, and it has now become routine to grow films in vacuum and transfer them immediately to an ARPES chamber after growth. Indeed, many ARPES systems have been developed that incorporate thin-film growth capabilities, such as those for molecular beam epitaxy and pulsed laser deposition; numerous achievements, such as the characterization of the electronic structure in high-critical-temperature (high-T) superconducting FeSe films\(^{110-111}\) and the metallic surface and interface states in oxide insulators\(^8\), have been enabled by these combined systems.

Another powerful way to prepare atomically clean and flat surfaces is the polish–sputter–anneal method, in which samples are polished, repeatedly sputtered and then annealed\(^{112}\). This method is suitable for certain materials that have strong chemical bonding, the surfaces of which can crystallize via annealing in vacuum. Recently, the polish–sputter–anneal method was applied to the study of the CoSi family of 3D topological semimetals\(^{113,114}\) and enabled the observation of multiple types of unconventional chiral fermions and helical Fermi arc surface states on the (001) surface.

For VUV ARPES, there are three types of light sources that have been successfully used: noble-gas discharge lamps\(^{115}\), synchrotron radiation\(^{116-118}\) and lasers\(^{119-121}\). At present, these three VUV light sources are complementary for ARPES experiments; their main features are summarized in Table 1 and are discussed in the following sections.

**Lamp-based VUV ARPES**

As the earliest photoemission light sources, noble-gas discharge lamps have been used widely in laboratory-based ARPES systems. These lamps use gas resonance lines as monochromatic light sources. A common example is the modern helium lamp, in which the radiation is generated through the microwave plasma technique\(^{118}\). Such a helium lamp mainly provides two discrete photon energies, 21.2 eV (He I\(\alpha\)) and 40.8 eV (He II\(\alpha\)), which can be selected by adjusting the angle of a diffraction grating in a monochromator. Additionally, the monochromator refocuses the beam to a spot size of ~1 mm with a typical flux of ~10\(^{12}\) photons per second.
Coulombic repulsion. photoelectrons induced by energy and momentum of the spectral redistribution of the Technical Reviews | OCTOBER 2019 | VOLUME 1 

The spectral redistribution of the inelastic mean free path of photoelectrons as a function of kinetic energy in nm. 

Fig. 2 | Synchrotron-based VUV ARPES in the study of TaAs surface states. Synchrotron-based vacuum ultraviolet (VUV) angle-resolved photoemission spectroscopy (ARPES) is a surface-sensitive technique that has been used to probe the surface states of Weyl semimetals, leading to the identification of Fermi arcs on the (001) surface of TaAs. a | Plot showing the universal curve of the inelastic mean free path of photoelectrons as a function of kinetic energy (bottom axis) and incident photon energy (top axis; calculated on the assumption that materials have a typical work function of 4 eV and electrons are located at the Fermi level). Coloured regions correspond to the typical photon energy ranges of VUV and soft-X-ray light. b | Schematic of a Weyl semimetal with spin-polarized Fermi arcs on its surfaces. The Fermi arcs connect the projections of two Weyl nodes of opposite chirality (indicated by the red and yellow points). The red arrows on the surfaces indicate the spin texture of the Fermi arcs. c | ARPES intensity plot of TaAs at the Fermi level ($E_F$) along the $\Gamma$–$\Gamma$ direction as a function of photon energy ($h\nu$). The observed Fermi surfaces (which appear as vertical lines, such as those indicated by the arrows) are constant despite variations in the photon energy, indicating that they are surface states. d | ARPES intensity plot at $E_F$ recorded on the (001) surface of TaAs ($k_y$ and $k_z$ are the momenta in the $y$ and $x$ directions, respectively). The yellow and red dots indicate the projection of Weyl points $W_1$ of opposite chirality. In panels d and e, $a_1$–$a_5$ are five different Fermi surfaces. e | Calculations of the surface states of TaAs at $E_F$ around Weyl points $W_1$ of opposite chirality (indicated by dashed circles as the chemical potential is slightly away from the nodes) along the $\Gamma$–$\Gamma$ direction. Panel a is adapted with permission from REF. [91], Wiley-VCH. Panel b is adapted with permission from REF. [59], Springer Nature Limited. Panel e is adapted from REF. [59], CC-BY-3.0.

In comparison to the other two VUV sources, discharge lamps are compact, and the beam position and intensity are stable. More importantly, the photon energy provided by discharge lamps is suitable for typical ARPES experiments, owing to reasonable photoemission cross sections $^{30,31,35} (~2$ MeV per atom for Au with He Ia light) and a negligible space-charge effect $^{12b,12c}$. However, discharge lamps have several limitations: relatively low photon flux (especially for high-resolution measurements) compared with a 7 eV laser $^{21,22}$, non-tunable photon energy with fixed or no polarization, a requirement of very flat sample surfaces owing to the large beam spot size, and leakage of noble gas into the ARPES chamber during measurements. Nevertheless, noble-gas discharge lamps are still the most popular and favourable laboratory-based light sources, and are particularly useful in the study of layered materials and thin films $^{311,313}$, which usually have large flat surfaces that can accommodate the large photon beam spot.

Synchrotron-based VUV ARPES

When electrons travelling at relativistic velocities are bent by a static magnetic field, strong electromagnetic radiation is emitted owing to the radial component of acceleration. This synchrotron radiation light can be applied to ARPES once correctly monochromatized. Indeed, benefiting from the continued development of synchrotron technology, especially the advent of third-generation synchrotron light sources, many synchrotron-based ARPES end-stations have been built and have gradually become the most powerful ARPES systems over the past 30 years $^{12–18}$.

Synchrotron-based ARPES has many notable advantages. First, by using a variable polarization undulator and high-resolution monochromator, the photon energy and polarization (linear or circular) of the beam can be readily tuned. The continuously tunable photon energy makes it possible to map the electronic structure in the entire 3D momentum space and to distinguish the
Laser Discrete values or discretely tunable within a limited range of several eV to tens of eV $<1$ (REFS 3,12,21) $\sim 10^5$–$10^6$ $\sim 100 \mu$m Tunable First to second Brillouin zones

Synchrotron Continuously tunable from several eV to thousands of eV $1$–$30$ (REFS 13,39) $\sim 10^{12}$–$10^{13}$ Tens of nanometres to hundreds of micrometres Tunable Several Brillouin zones

Discharge lamp Discrete values between several eV and tens of eV$^b$ $\sim 1^b$ $\sim 10^{12}$ $\sim 1$ mm Non-tunable First to second Brillouin zones

ARPES, angle-resolved photoemission spectroscopy. aExamples of common discharge lamps include the He lamp (with photon energies of 21.2 eV for He I and 40.8 eV for He II) and the Xe lamp (with a photon energy of 8.4 eV for Xe I). bData for the VUVk He lamp.

Application to topological Weyl semimetals. The continuously tunable VUV light with which most synchrotron-based ARPES sets-ups are equipped has enabled many important breakthroughs in the study of topological materials, such as the observation of surface Fermi arcs in the Weyl semimetal TaAs (REFS 33,34,37,88–95). Weyl semimetals are a class of materials that can be regarded as 3D analogues of graphene wherein the bulk non-degenerate electronic bands disperse linearly along all momentum directions through a node near $E_F$, called a Weyl node, which can be viewed as a singular point of Berry curvature or a ‘magnetic monopole’ in momentum space. Weyl nodes can appear in pairs of opposite chirality in a real material if the spin-doublet degeneracy of the bands is removed, which happens when either time-reversal or inversion symmetry is broken. The low-energy excitations near the Weyl nodes behave like Weyl fermions, which were originally proposed by Hermann Weyl in 1929 as massless solutions of the Dirac equation. However, Weyl fermions have not yet been observed in particle physics; hence, the physical realization of Weyl fermions in Weyl semimetals becomes more notable.

One hallmark of a Weyl semimetal is the existence of pairs of Fermi arcs at two opposite crystal surfaces. Each Fermi arc connects the projections of two opposite-chirality bulk Weyl nodes in the surface Brillouin zone (Fig. 2b). Synchrotron-based VUV ARPES, with its unique tunable light, has played a central part in identifying surface Fermi arcs in the TaAs family of semimetals (TaAs, TaP, NbAs and NbP). As an example, on performing photon-energy-dependent measurements along the $\bar{F}–\bar{Y}$ direction in TaAs (Fig. 2c), the surface states at $k_y = 0.4 \AA^{-1}$ (where $k_y$ is the momentum in the $y$ direction) could be clearly identified as they do not show any noticeable change in $k_y$ upon variation of the incident photon energy. To identify the surface Fermi arcs, it is necessary to subsequently map out the Fermi surfaces by collecting photoelectrons of different emission angles. Indeed, a surface Fermi arc (labelled $a$ in Fig. 2b) is clearly resolved in the Fermi surface map of TaAs and shows very good agreement with the calculated surface states on the (001) As-terminated surface (Fig. 2e).

Laser-based VUV ARPES

On the basis of their high photon flux and narrow bandwidth, lasers are a natural choice for laboratory-based ARPES. However, the photon energies generated by lasers are usually too low to overcome the 4–5 eV work function of metallic materials in an ARPES measurement. Great efforts have been made in recent years to develop high-energy (>5 eV) lasers that are suitable for ARPES, and by taking advantage of nonlinear optical processes, long-sought VUV lasers have been achieved.

The VUV laser light required for the photoemission process is typically generated in one of two ways. One way is to exploit the process of harmonic generation through nonlinear optical crystals. The most commonly used nonlinear optical crystals are BaB$_2$O$_4$ and KBe$_2$BO$_3$F$_2$. Using BaB$_2$O$_4$ crystals, a Ti:sapphire laser (with a photon wavelength and energy of ~800 nm and 1.55 eV, respectively) with a sub-picosecond pulse width can output ~6 eV laser light, whereas with KBe$_2$BO$_3$F$_2$ crystals, a Nd:YVO$_3$ laser (~1.064 nm; 1.17 eV) with a femtosecond pulse width can produce ~7 eV laser light with an ultranarrow bandwidth and a high flux. Alternatively, VUV laser light can be generated using high harmonic generation (HHG) or multiphoton excitations in noble gases; this approach generates laser light with a higher photon energy (approximately tens of eV) than that of nonlinear optical crystals (<10 eV). For example, a Yb-doped fibre laser,
when coupled with a xenon gas cell, can output ~11 eV laser light\(^{11}\), and a Ti:sapphire laser with a gas cell can generate discretely tunable laser light in the range 15–40 eV with attosecond\(^{16}\) or femtosecond pulse durations\(^{17}\). However, these sources usually have a relatively low photon flux (≤10\(^{13}\) photons per second) compared with 6–7 eV lasers, owing to the low generation efficiency of the HHG process and unavoidable loss of flux in the optics after generation. Moreover, the femtosecond or attosecond light pulses usually have a low energy resolution (ΔΕ, of the order of tens of meV), which is constrained by the Fourier-transform limit. This limit, when expressed in terms of convenient units for Gaussian pulses with a pulse width Δτ, is given by\(^{20}\)

\[ ΔτΔΕ \geq 1,825 \text{ fs meV} \]  

(5)

where Δτ and ΔΕ are both quoted at full-width at half maximum. Although HHG laser light is not suitable for high-resolution ARPES measurements owing to the low energy resolution, the short pulse widths make it possible to perform time-resolved ARPES measurements at relatively high photon energies (discussed further below in the section on time-resolved ARPES).

Laser-based ARPES has many advantages. One major benefit of laser ARPES with low-energy (~6–7 eV) photons is a substantial gain in the resolution of the in-plane momentum. From Eq. 2,

\[ Δk_∥ \propto \sqrt{2mE_{\text{kin}}/h^2 \cdot \cosθ \cdot Δθ} \]  

(6)

Therefore, for the same θ and Δθ, photoelectrons with a low kinetic energy result in a better in-plane momentum resolution. Moreover, the high photon flux (~10\(^{15}\) photons per second) and extremely narrow bandwidth (<1 meV) of some 6–7 eV laser light sources with long pulse widths (>2 ps) also make it possible to perform ARPES measurements with a very high energy resolution and with high data acquisition efficiency. Thus, the low-energy laser light sources have the advantage of ultrahigh energy and momentum resolution. Indeed, laser-based ARPES systems with a photon energy of ~7 eV and energy resolution better than 1 meV have been achieved by using harmonic generation in a KBe\(_2\)BO\(_3\)F\(_2\) nonlinear crystal\(^{21,22}\). In addition, unlike noble-gas discharge lamps, it is easy to control the polarization of the laser light and to tune the beam spot size to the sub-100 μm range. More importantly, pulsed laser light introduces a new degree of freedom to ARPES — time resolution — which we discuss in more detail below.

As with the other VUV light sources, laser light has its limitations. The main disadvantages are the limited tunability of the photon energy, low photoemission cross sections for some materials (for example, the cross section is <0.01 Mb per atom for Au with a 10 eV laser) and the inability to access electrons far from the centre of the Brillouin zone with a photon energy of 6–7 eV. To some extent, the application of HHG laser light with discretely tunable photon energies of tens of eV has overcome these problems. However, it remains difficult to obtain high-stability and high-flux photons with a narrow bandwidth.

**Application to Fe\(_{0.55}\)Te\(_{0.45}\)** To demonstrate the high-resolution power of laser-based ARPES, we highlight one application: the observation of a topological surface Dirac cone on the (001) surface of Fe\(_{0.55}\)Te\(_{0.45}\)\(^{15}\). As the simplest iron-based superconductor (Fig. 3a), Fe(FeSe)\(_{0.45}\)\(_{15}\) has attracted tremendous attention since its discovery in 2008\(^{17}\). More recently, calculations\(^{122,124}\) have predicted the existence of topological order in Fe\(_{0.55}\)Te\(_{0.45}\). The substitution of Se with Te lowers the energy of the p\(_{y}\) band and leads to band inversion in Fe\(_{x}\)Te\(_{1−x}\)Se\(_{0.45}\) between the p\(_{x}\)(−) and d\(_{x−y}\)(−) or d\(_{x−y}\)(+)+ bands along the Γ–Z direction (Fig. 3b), where + and − denote even and odd parity under spatial inversion, respectively. Eventually, owing to the opening of the spin–orbit coupling gap between these two inverted bands, this material naturally hosts strong topological surface states inside the gap at the Γ point of the (001) surface (Fig. 3c). The limited energy and momentum resolution of ARPES has, however, made it difficult to clearly resolve these surface states\(^{15,136}\). Only in 2018 were the Dirac surface states revealed through the use of high-resolution laser ARPES\(^{139}\). With this technique, a Dirac-cone surface state near E\(_F\) was clearly resolved from a high-resolution curvature intensity plot (Fig. 3d). More importantly, when Fe\(_{0.55}\)Se\(_{0.45}\) enters the superconducting state, superconductivity is also induced in the topological surface states (Fig. 3e), and temperature-dependent and momentum-dependent measurements of the induced superconducting gap were recorded (Fig. 3f). The energy distribution curves at slightly higher temperatures show evidence of the hole branch of Bogoliubov quasiparticles (the shoulders above E\(_F\) in Fig. 3f), demonstrating the superconducting nature of the topological surface state. In addition, the superconducting gaps on the topological surface state are isotropic in momentum (Fig. 5g,h), which indicates that the surface state has an s-wave superconducting pairing symmetry. Therefore, Fe\(_{0.55}\)Se\(_{0.45}\) provides an excellent platform for realizing Majorana zero modes based on the Fu–Kane model\(^{140}\), in which a single spin–momentum locked Dirac surface state with proximitized s-wave superconducting pairing can have the same role as a p-wave superconducting order parameter. Indeed, the Majorana zero modes have recently been observed in the superconducting vortex cores of FeTe\(_{0.55}\)Se\(_{0.45}\) with scanning tunnelling microscopy\(^{138}\).

**Bulk-sensitive soft-X-ray ARPES** Synchrotron-based VUV ARPES, usually with photon energies ≥20 eV, is an ideal tool for studying the surface electronic structure of many materials, especially topological materials. However, it is difficult to probe the bulk states of 3D materials with ≥20 eV VUV light owing to the extremely short escape depths of the resulting photoelectrons (Fig. 2a) and the ill-defined k\(_z\). Although the photoelectrons ejected by 6–7 eV lasers have a longer escape depth (Fig. 2a) compared with those of ≥20 eV VUV sources, the non-tunable photon energy and stronger final-state effects\(^{49}\) due to low-speed photoelectrons make it nearly impossible to probe the intrinsic 3D band structure of bulk states.

The only reliable way to examine 3D bulk states, especially the band dispersion along the k\(_z\) direction, is to use...
The high energy and momentum resolution of laser-based vacuum ultraviolet (VUV) angle-resolved photoemission spectroscopy (ARPES) has enabled the identification of the Dirac surface states in FeTe$_{0.55}$Se$_{0.45}$. a | Crystal structure (left) and the bulk and projected (001) surface Brillouin zones (right) of Fe(Fe,Se). b | Calculated band structure of FeTe$_{0.55}$Se$_{0.45}$ along the Γ–M and Γ–Z directions (where $t$ = 100 meV). The dashed box shows the spin–orbit coupling gap of the inverted t$_2$g–e$_g$ bands. c, d | Calculated band structure (panel c; t = 100 meV) and curvature intensity plot of ARPES data (panel d) along the Γ–M direction (where $k$ is the momentum in the Γ–M direction). The topological surface states (TSSs) connecting the bulk valence band (BVB) and bulk conduction band (BCB) are clearly resolved in both images. The ARPES data in panel d were recorded with a p-polarized 7 eV laser at 15 K. e | Schematic of the bulk and surface superconducting (SC) states in FeTe$_{0.55}$Se$_{0.45}$. Below the superconducting transition temperature ($T_c$), the bulk states open s-wave SC gaps (where $\Delta$ is the SC gap size); these bulk states are topologically trivial because of their spin degeneracy (black curves). Owing to the superconducting proximity effect, the TSS also opens an isotropic gap at temperatures ($T < T_c$) and is topologically superconducting (TSC) as a consequence of the spin polarization (blue and red curves). f | Raw energy–distribution curves measured at different temperatures for a $k$ point on the Fermi surface. The shoulders above the Fermi level ($E_F$) signify SC Bogoliubov quasiparticles. The energy positions of the coherence peaks in the energy-distribution curves correspond to the SC gap size. g | Symmetrized energy–distribution curves of the Dirac surface states at different Fermi wavevectors (indicated in panel h) recorded at $T = 2.4$ K. h | Polar representation of the measured SC gap size in panel g. The measured SC gaps (solid markers) at different polar angles ($\phi$) almost have the same value, demonstrating that the SC gaps of surface states are also isotropic in momentum. The hollow markers are a mirror reflection of the solid markers, and the vertical lines passing through the solid or hollow markers indicate the error bars. $E$, energy. Adapted with permission from REF$^{156,16}$, AAAS.

There are three main reasons that soft-X-ray ARPES facilitates the study of bulk states. First, the higher kinetic energy of the photoelectrons ejected by soft X-rays results in an increase in the photoelectron escape depth by a factor of 2–4 compared with that of 20 eV VUV light (for which the escape depth is ~5 Å); this increase greatly enhances the sensitivity to the bulk electronic states. Second, from the Heisenberg uncertainty principle, the momentum broadening of photoelectrons along the $k_z$ direction, $\Delta k_z$, is given by $\Delta k_z \propto d^{-1}$, where $d$ is the photoelectron escape depth. Thus, an increase in the escape depth with soft X-rays leads to a decrease of $\Delta k_z$ and an improvement in the $k_z$ resolution, which enables accurate investigations of the bulk states of 3D materials in the entire momentum space. Third, in the photoemission process with high-energy soft-X-ray incident light, the final states are truly free-electron-like, which enables precise determination of the $k_z$ value from Eq. 5.

In addition to the high sensitivity to bulk electronic states, soft-X-ray ARPES has other advantages, such as simplified matrix elements and, compared with VUV ARPES, a lower sensitivity to the sample surface quality. However, soft-X-ray ARPES also suffers from several problems. The main difficulty is a decrease in the valence-band cross section by two to three orders of magnitude compared with the VUV energy range$^{119}$. This decrease is due to reduced wavefunction overlap between the spatially rapidly varying final states and the spatially smooth valence states. Thus, only initial electronic states near a very small region around the ion core, the wavevector of which matches the large wavevector of the high-energy final states, contribute significantly
to the photoemission intensity\textsuperscript{13,14}. This signal loss can be compensated by a high flux of incident photons. For example, owing to the advances in synchrotron radiation sources and beamline instrumentation, the soft-X-ray ARPES end-station at the Swiss Light Source has successfully compensated the signal loss with a soft-X-ray flux of >10\textsuperscript{13} photons per second\textsuperscript{14}.

Another difficulty with soft-X-ray ARPES, in sharp contrast to low-energy laser ARPES, is that the use of high-energy soft-X-rays leads to a reduction in the in-plane momentum resolution (Eq. 6). Therefore, spectrometers need to be equipped with a higher angular resolution to improve the \( k_z \) resolution. Furthermore, compared with VUV ARPES, soft-X-ray ARPES also suffers from a loss of energy resolution. More specifically, the energy resolution of VUV ARPES can be better than 1 meV, whereas for soft-X-ray ARPES, the energy resolution varies from tens of meV to 100 meV depending on the photon energy. Finally, we caution that the photon momentum of soft X-rays (~0.25 Å\textsuperscript{−1} at 500 eV) may not be negligible compared with the typical size of the Brillouin zone (~1–2.5 Å\textsuperscript{−1}). Hence, the momentum conservation law in Eq. 2 may no longer hold.

**Application to topological semimetals**

To illustrate the study of bulk electronic structure using soft-X-ray ARPES, and to demonstrate how central the improvement in \( k_z \) resolution has been in this regard, we use 3D topological semimetals TaAs (REFS\textsuperscript{33,35}) and MoP (REF\textsuperscript{35}) as examples. Topological semimetals, which have symmetry-protected band-crossing points, have become one of the most intensively studied fields in condensed matter physics. The most famous examples are Dirac\textsuperscript{55}–\textsuperscript{57} and Weyl semimetals\textsuperscript{55,34,37,58–59}, in which two doubly or singly degenerate bands cross, forming four-fold Dirac points or two-fold Weyl points, respectively (FIG. 4a). New types of crystal symmetry-protected band crossings have been theoretically predicted in condensed matter systems, and the corresponding low-energy excitations have no high-energy counterparts\textsuperscript{141–149}. These excitations, also referred to as unconventional fermions, could challenge our existing wisdom...
about the classification and properties of fermions. For example, topological semimetals with three-fold band crossings have been predicted in several materials with WC-type structures\textsuperscript{144–146}. The quasiparticle excitations near the band-crossing points are three-component fermions, which can be viewed as the intermediate species between the four-component Dirac and two-component Weyl fermions.

With the ability to probe more deeply into a sample, soft-X-ray ARPES has had a key role in the detection of bulk Weyl points in TaAs\textsuperscript{145,146} and triply degenerate points in MoP\textsuperscript{15} and WC\textsuperscript{16}. For the Weyl semimetal TaAs, first-principles calculations indicated that there are 12 pairs of Weyl points in the bulk Brillouin zone\textsuperscript{16}. However, owing to the short escape depth of the photoelectrons excited by VUV light, the bulk Weyl points could not be resolved in VUV ARPES experiments\textsuperscript{16} (fig. 2d). However, with soft-X-ray ARPES, the bulk Weyl bands are clearly resolved\textsuperscript{13,34,150}. The measured electronic states in the $\Gamma$–$\Sigma$–Z–$\Sigma$ plane, which is perpendicular to the cleaved sample surface, clearly exhibit a periodic modulation upon varying the incident soft-X-ray photon energy (hence varying $k_z$; fig. 4b), confirming the bulk nature of the detected spectra. The bulk Weyl points (labelled W1 in fig. 4c) were confirmed from the measured M-shaped band dispersion, with the two peaks corresponding to a pair of Weyl points of opposite chirality.

The benefits of soft-X-ray ARPES are further illustrated in studies on MoP\textsuperscript{15}. Owing to its 3D crystal structure, the top layer of the cleaved surface of MoP is amorphous (fig. 4d). Therefore, with VUV ARPES, no obvious Fermi surface is observed at 60 eV (fig. 4d) owing to a combination of the limited detection depth of VUV ARPES and the angle-smearing effect of scattering in the amorphous layer. However, upon increasing the photon energy to 453 eV, the bulk states are clearly seen (fig. 4f), enabling the detection of the predicted triply degenerate point. Indeed, by using high-precision measurements of the band dispersions, it is possible to resolve the triply degenerate point with soft X-rays (fig. 4g), and this is well reproduced by calculations (fig. 4h).

**Spin-resolved ARPES**

Historically, the development of ARPES has been driven by the demands of research for a specific class of materials. For example, the discovery of the high-$T_c$ superconductors promoted the development of high-resolution laser ARPES\textsuperscript{21,22}. Similarly, the recent discovery of topological insulators and non-centrosymmetric materials has stimulated the development of high-performance spin-resolved ARPES\textsuperscript{16–42}.

The integration of spin detectors into ARPES spectrometers enables spin-resolved ARPES, and substantial efforts have been made to develop compact spin polarimeters. Examples of detectors include those using spin-polarized low-energy electron diffraction\textsuperscript{15}, diffuse scattering\textsuperscript{16}, Mott scattering\textsuperscript{17} and very low-energy electron diffraction (VLEED)\textsuperscript{18}. Most of these spin detectors are based on the asymmetry ($A_m$) of preferential spin scattering, which can be written as $A_m = (I_m - I) / (I_m + I)$, where $I_m$ and $I$ are the partial intensities of electrons with magnetic moments parallel or antiparallel, respectively, to the target magnetization direction of the detector. The final spin polarization ($P$) is proportional to the asymmetry: $P = A_m / S$, where $S$ is the spin sensitivity (a coefficient called the Sherman function in Mott detectors), which can be determined by measuring a fully polarized electron beam with $A_m = 1$. The most widely used spin polarimeters are Mott detectors\textsuperscript{19} and VLEED detectors\textsuperscript{15}, which use heavy elements (such as Au and Th) and ferromagnetic thin films (such as Fe(001) and Co(001) films) as scattering targets, respectively. A Mott detector uses the scattering asymmetry induced by the spin–orbit interaction when high-energy ($\sim 20–100$ keV) electrons scatter off a heavy-element target. By contrast, a VLEED detector takes advantage of the exchange scattering asymmetry of very low-energy (<10 eV) electrons with ferromagnetic thin films. Mott detectors have the advantage of using scattering targets with long lifetimes, but the scattering rate ($\sim 10^{-2}$) is very low owing to the low cross sections of the high-kinetic-energy electrons. By contrast, VLEED detectors have a much higher efficiency (approximately two orders of magnitude higher than that of Mott detectors), owing to the higher scattering probability of very low-energy electrons and improved spin sensitivity. However, the ferromagnetic thin-film target has to be regenerated frequently, typically once every few weeks, because of degradation of the thin film.

Spin-resolved ARPES is one of the most powerful techniques for directly measuring the spin texture of electronic states. However, it also has several shortcomings. First, the energy resolution and angular resolution of the detectors still lag behind those of non-spin-resolved electron analysers. Second, the detectors have a low efficiency, which is due to the small scattering cross section and the single detection channel, in which only photoelectrons emitted within a small solid angle are recorded. Notably, recent advances in spin detectors\textsuperscript{23–26} and light sources have greatly improved the performance of spin-resolved ARPES. For example, combining a state-of-the-art spin detector (a Scienta Omicron DA30-L spin spectrometer; fig. 5a) and a high-photon-flux VUV laser, a high efficiency spin-resolved ARPES system with an energy resolution of 1.7 meV was achieved\textsuperscript{40}. The DA30-L spin spectrometer combines a hemispherical electron analyser with twin VLEED spin detectors (VLEED-B and VLEED-W). The electron analyser is equipped with electron deflectors, whereby an applied electric field is used to select photoelectrons of desired emission angles. The DA30-L analyser therefore enables detailed $k$-space mapping of 2D in-plane electronic structure, $E(k_{x,y})$, at a fixed sample geometry with acceptance angles $\theta_x = 30^\circ$ and $\theta_y = 24^\circ$, where $x$ and $y$ are the directions parallel and perpendicular to the entrance slit, respectively. The twin VLEED spin detectors use oxidized iron films as scattering targets, and the film target of VLEED-B (VLEED-W) is magnetized by solenoid coils in the $x$ and $z$ ($y$ and $z$) directions, which correspond to the spin-polarization directions of $P_x$ and $P_y$ ($P_y$ and $P_z$) on the sample axis. Thus, the DA30-L spin spectrometer enables analysis of the 3D spin vector of photoelectrons emitted in the acceptance cone.
of \((\theta \times \theta) = (30^\circ \times 24^\circ)\) without sample rotation. Note, however, that, in general, the spin-polarization signal can be complicated by matrix element effects\(^{161-163}\). Thus, photon-energy-dependent and photon-polarization-dependent spin-resolved ARPES measurements are needed to check whether the spin signal is intrinsic to the magnetic moment of the photoelectrons\(^{43,46}\).

**Application to 3D topological insulators**

The discovery of 3D topological insulators has helped to drive the development of spin-resolved ARPES. A topological insulator is a state of quantum matter that features an energy gap in the bulk but gapless Dirac-cone surface states that reside inside the bulk insulating gap\(^{4,44,22,25-27}\). The distinctive feature of a Dirac-cone surface state is the spin–momentum locking pattern\(^{4,44}\), which manifests as a spin texture that winds in a circle around a constant-energy contour of the Dirac-cone surface state (FIG. 5a). To highlight the role of spin-resolved ARPES in the examination of the spin texture of these surface states, we use Bi\(_2\)Te\(_3\) as an example because it hosts a clean Dirac cone near \(E\_F\) (FIG. 5c). Indeed, the spin–momentum locking feature of the surface electrons in Bi\(_2\)Te\(_3\) was observed with spin-resolved ARPES\(^{43}\). The measured spin-polarization spectra along \(\Gamma - \text{M}\) in the \(x, y\) and \(z\) directions (FIG. 5d,e) show that there is no clear spin polarization in the \(x\) and \(z\) directions within the experimental resolution, whereas clear polarization signals of equal magnitude and opposite signs are observed in the \(y\) direction. This finding implies that the surface electrons of opposite momenta also have opposite spin textures, confirming the spin–momentum-locking scenario. Besides spin-resolved ARPES, the spin textures of electronic states can also be extracted from the spin-dependent differential absorption of left-circularly versus right-circularly polarized light, which is the basis of circular dichroism\(^{164-166}\).

**Ultrafast time-resolved ARPES**

Non-time-resolved ARPES serves as an excellent tool for probing the band dispersion of equilibrium states. However, it remains challenging to distinguish and quantify many-body interactions (for example, electron–electron and electron–phonon interactions) of correlated materials. Generally, electron–electron, electron–spin and electron–phonon interactions occur...
on timescales of femtoseconds, tens of femtoseconds and picoseconds respectively. Therefore, the use of sub-picosecond or even sub-femtosecond laser pulses in pump–probe experiments enables the coupled interactions between the charge, spin, lattice and orbital degrees of freedom to be disentangled. In particular, the combination of pump–probe optical spectroscopy with ARPES, namely time-resolved ARPES, provides direct insight into the energy and momentum dependence of these ultrafast dynamics.

In a time-resolved ARPES experiment, a femtosecond laser pulse (the ‘pump’) is used to perturb a material into a non-equilibrium state. Subsequently, a second time-delayed pulse (the ‘probe’), which is typically in the VUV range, is used to excite photoelectrons out of the sample; these photoelectrons are then analysed by an electron spectrometer. By varying the time delay between the pump and probe pulses, it is possible to obtain insight into the time-dependent processes involved in the excitation and relaxation of the transient states. Besides providing energy-resolved and momentum-resolved transient spectra in the time domain, time-resolved ARPES also allows for the investigation of unoccupied states above $E_F$ by populating them with photoexcited electrons while keeping the band structure minimally disturbed. It is worth mentioning that besides time-resolved ARPES, the unoccupied states can also be detected using two-photon photoemission spectroscopy (such as the RoentDek delay-line detectors) to overcome this limitation by enabling multiple electrons to be detected simultaneously per pulse. The relatively high efficiency of ARTOF analysers makes them ideal detectors for pulsed lasers. However, the inherent requirement for the pulsed beams to have repetition rates < 3 MHz, for compatibility with the delay-line detectors, prohibits the application of ARTOF analysers to light sources that have repetition rates > 3 MHz, including most synchrotron light sources, gas-discharge lamps and quasi-continuous lasers. Another drawback of ARTOF analysers is the high background noise of scattered photons off the sample surface, which travel in a straight line to the delay-line detector. By contrast, scattered photons are not an issue for hemispherical detectors because the scattered photons are either blocked by the narrow entrance slit at one end of the hemisphere or stopped in the curved hemisphere if a small fraction passes the slit, and thus do not reach the detector at the opposite end of the hemisphere.

The rapid development of commercially available, high-repetition-rate (>10 kHz), amplified femtosecond lasers has led to the increased use of time-resolved ARPES. In the past 10 years, various time-resolved ARPES systems with different energy resolution (tens to hundreds of meV), temporal resolution (approximately hundreds of attoseconds to hundreds of femtoseconds) and pump frequencies (approximately tens of meV to several eV) have been developed, enabling the electron dynamics in correlated materials to be probed. Moreover, by carefully choosing the pump frequency, it is possible to observe new quantum states induced by light, such as Floquet–Bloch states discussed below.

### Application to grey As and Bi$_2$Se$_3$

The capability of time-resolved ARPES as a tool to directly probe the dynamics of transient states and unoccupied states is clearly illustrated in the study of grey As. Grey As exhibits non-trivial Rashba–Spin–Luttinger liquid states on the (111) surface; these non-trivial states arise from a bulk band inversion caused by the crystal field. To elucidate the non-trivial band topology of these Shockley states, time-resolved ARPES measurements were performed on the (111) surface to probe the unoccupied states above $E_F$. The ARPES
**Fig. 6 | Time-resolved ARPES studies of grey As and Bi₂Se₃.** Time-resolved angle-resolved photoemission spectroscopy (ARPES) can be used to probe the dynamics of transient states and to manipulate the electronic structure of a material. 

**a** | Schematic of a time-resolved ARPES system with an angle-resolved time-of-flight (ARTOF) spectrometer. In this set-up, a pump laser with photon energy \( h\nu_p \) is used to excite a material into a non-equilibrium state. Subsequently, a probe laser pulse with photon energy \( h\nu_p \) arriving after a time delay \( \Delta t \), is used to eject photoelectrons from the sample. The photoelectrons are then analysed by an ARTOF detector, which calculates the electron kinetic energy by measuring the flight time from the sample to the detector and deduces the electron momentum from the strike position on the 2D delay-line detector (DLD). An ARTOF analyser can therefore measure the entire band structure \( E(k_x, k_y) \) (where \( E \) is the energy, and \( k_x \) and \( k_y \) are the momenta in the \( x \) and \( y \) directions, respectively) simultaneously; an example is shown for the surface Dirac cone in Bi₂Se₃ (top left). 

**b** | Measured band dispersions of grey As on the (111) surface above and below the Fermi level \( E_F \) along the \( \Gamma \rightarrow K \) direction. The top parts show the band dispersions above \( E_F \), recorded using high-resolution static laser ARPES measurements. The bottom parts show the band dispersions below \( E_F \), recorded using high-resolution static laser ARPES measurements. 

**c** | Slab calculations of surface bands with opposite spin orientations (shown in red and blue) of the (111) surface of grey As and the bulk bands (indicated in black) along the \( \Gamma \rightarrow K \) direction. The double-headed red arrow indicates the linear polarization \((\Gamma \rightarrow K)\) of the pump laser projected onto the sample surface. 

**d** | Time-resolved ARPES intensity snapshots of grey As along the \( \Gamma \rightarrow K \) direction with various pump–probe delay times. The snapshots show the excitation or decay process of the electronic states above \( E_F \). 

**e** | Schematic of an ARTOF spectrometer, showing the flight time from sample to detector and deduces the electron momentum from the strike position on the 2D delay-line detector (DLD). An ARTOF analyser can therefore measure the entire band structure \( E(k_x, k_y) \) (where \( E \) is the energy, and \( k_x \) and \( k_y \) are the momenta in the \( x \) and \( y \) directions, respectively). 

**f** | Illustration showing the experimental geometry of a time-resolved ARPES measurement with a \( p \)-polarized pump laser. The cone structure represents the surface Dirac cone (centred at the \( \Gamma \) point) of Bi₂Se₃. The double-headed red arrow indicates the linear polarization \((\Gamma \rightarrow K)\) of the pump laser, and the blue arrow denotes the electric field \((E_F)\) of the pump laser projected onto the sample surface. 

**g** | Time-resolved ARPES intensity plots of Bi₂Se₃ along the \( k_y \) direction at \( \Delta t = -0.5 \) ps with a \( p \)-polarized pump laser. 

**h** | Time-resolved ARPES intensity plots of Bi₂Se₃ along the \( k_y \) (panel g) and \( k_x \) (panel h) directions at \( \Delta t = 0 \) with a \( p \)-polarized mid-infrared pump laser. The red arrows in panel h indicate the avoided crossing gaps. 

**i** | Sketch of the ‘dressed’ replica bands of different orders induced by the mid-infrared pump pulse. Avoided crossings occur along the \( k_y \) direction, leading to a bandgap of 2D. BCB, bulk conduction band; BVB, bulk valence band; \( E \), energy; \( h\nu \), photon energy of pump pulse; SS, surface state. Panel a is adapted with permission from ref. 166, APS. Panels b–d are adapted with permission from ref. 166, APS. Panels f–h are adapted from ref. 166, Springer Nature Limited. Panels e and i are with permission from ref. 166, AAAS.
data (FIG. 6b) clearly reveal a pair of parabolic bands, which split along both the Γ−M and Γ−K directions but are degenerate at the Τ point. In the Γ−K direction, one band disperses into the conduction band, while the other turns back and merges into the valence band. This behaviour shows excellent agreement with the calculated band structure (FIG. 6c), thus providing strong evidence for the non-trivial band topology of the Shockley states. Furthermore, by varying the pump-probe delay time, the dynamics of these unoccupied states, such as the characteristic relaxation time, can be revealed (FIG. 6d).

Time-resolved ARPES can also be used to detect new quantum states that emerge from the coherent interaction between electrons and photons. Among the most intriguing achievements has been the observation of Floquet–Bloch states in Bi2Se3. According to Bloch’s theorem, a spatially periodic potential in a lattice results in the replication of band dispersion in momentum, namely, Bloch states. In analogy to Bloch states, a temporally periodic electromagnetic field from an intense pump laser pulse leads to replicas of the bands in energy, known as Floquet states. Floquet–Bloch states are then a result of periodic potentials in both time and in space. Floquet–Bloch states were demonstrated on the surface of a Bi2Se3 sample by using a p-polarized pump pulse (FIG. 6e) with an energy lower than that of the bulk bandgap to avoid the excitations of abundant electron–hole pairs in the bulk, which can mask the observation63,182. Without pumping, a single Dirac-cone surface state resides inside the bulk bandgap (FIG. 6f). When pumping Bi2Se3 with p-polarized laser pulses, periodic duplicates of the Dirac bands begin to appear along the energy axis (FIG. 6g,h), separated by the pump photon energy. These dispersive band replicas cross upon moving away from the Τ point, giving rise to new crossing points (FIG. 6i). Notably, an energy gap of 2Δ opens along the kx direction at the expected crossings while these crossings remain gapless along the ky direction (FIG. 6i–j). This distinction between the two momentum directions is because the perturbing Hamiltonian associated with a p-polarized pulse commutes with the Dirac Hamiltonian for electrons with momentum along kx but not for those along ky (REFS182,183). This gap opening along the ky direction distinguishes the observed band replicas from those caused by laser-assisted photoemission, for which no gaps are expected64. The above observations provide direct evidence of the photon-dressed Floquet–Bloch states in solids and may pave the way for optical manipulation of new phases.

The capability of time-resolved ARPES goes far beyond the above two examples, and it has been widely used in studying the momentum-resolved electronic dynamics of many classes of materials, including high-Tc superconductors (for example, cuprates65,184-192), iron-based superconductors66,193-199, density wave systems60,62-68,170,178,199,4 topological insulators64,65,182,200-207, graphene108-211 and other strongly correlated systems212,213. For example, in addition to Floquet–Bloch states, the difference in dynamics of scattering between surface and bulk states has also been observed in Bi(I)Se(II)Te(III) (REFS206,207). Apart from incoherent transient dynamics, coherently excited bosonic modes have also been observed in numerous materials, such as transition metal dichalcogenides208,209, rare-Earth tritellurides210,216 and FeSe/SrTiO3 films216, often yielding important information about the electron–boson coupling in these systems.

Spatially resolved ARPES
Typical synchrotron-based ARPES systems have a spatial resolution of ~100 μm, depending on the spot size of the light source; thus, samples to be measured should have a flat surface with dimensions greater than 100 μm. However, many interesting materials or single-crystalline domains are smaller than 100 μm, such as heterostructures and microscale and nanoscale materials. In consideration of the above cases, spatially resolved ARPES with micrometre or sub-micrometre spatial resolution has been developed at several synchrotron light sources4,48-51,214.

Spatially resolved ARPES can be viewed as a combination of ARPES and scanning photoemission microscopy. Depending on the beam spot size, spatially resolved ARPES is referred to as micro-ARPES or nano-ARPES for microscale or nanoscale beams, respectively. To focus X-rays to a microparized or nanosized spot at the sample, advanced optics have been developed, including Schwarzschild optics4 and the Fresnel zone plate214. As an example, the nano-ARPES end-station at the MAESTRO beamline at the Advanced Light Source (California, USA) uses a Fresnel zone plate to focus the beam, followed by an order-sorting aperture to eliminate high-diffraction orders (FIG. 7a). In general, a high-precision sample stage is also needed to ensure precise nanometre scanning and positioning of the sample. The final spatial resolution is determined by the resolution of the X-ray optics and the mechanical and thermal stability of the sample stage.

The development of spatially resolved ARPES has made it possible to map out the band structures of materials with high spatial resolution while maintaining the advantages of ARPES, such as good energy and momentum resolution. However, the shortcomings of spatially resolved ARPES include low count rates owing to the low focusing efficiency (~1%) of the focusing optics, strong space-charge effects due to the nanosized beam spots, and limited photon–energy choices owing to the complexity of designing photon–energy-dependent focusing optics. A complementary technique, photoemission electron microscopy215, also enables spatially resolved electronic structure measurements. However, the corresponding energy and angular resolution (typically >100 meV and 1°, respectively) of this technique are usually worse than those of spatially resolved ARPES (which has an energy resolution of the order of tens of meV and an angular resolution of ~0.1°).

Application to weak topological insulators
Spatially resolved ARPES has been applied to the study of electronic structure of microscale and nanoscale materials and domains, including heterostructures22,23,216-221, SBs2Te3 nanowires224,225 and the weak topological insulator β-Bi2I3 (REF224). We highlight the
Spatially resolved ARPES in the study of the topological surface states of \(\beta\)-Bi\(_4\)I\(_4\). Spatially resolved angle-resolved photoemission spectroscopy (ARPES) is a tool for studying the electronic structure of nanoscale or microscale materials and domains, such as the topological surface states of the weak topological insulator (WTI) \(\beta\)-Bi\(_4\)I\(_4\). a | Schematic of a nano-ARPES system. A set of Fresnel zone plates (FZPs) is used to focus the incoming beam, followed by an order-sorting aperture (OSA) to suppress unwanted diffraction orders. b | Schematics of the topological surface states (TSSs) of \(\beta\)-Bi\(_4\)I\(_4\) in real space (left) and their band dispersions in \(k\)-space (right), where \(k_x\) and \(k_y\) are the momenta in the \(y\) and \(x\) directions, respectively. The small spheres in the left part represent the electrons, with the black arrows indicating the spin texture of the electrons in real space; the blue and red arrows indicate the direction of electron movement. The black arrows in the right part indicate the spin texture in \(k\)-space. c | Calculated Dirac band dispersions of TSSs at the \(\Gamma\) (left) and \(\bar{Z}\) (right) points of the side (100) surface. d | Schematic showing the preparation of the (100) side surface for nano-ARPES measurements. Fresh side (100) surfaces can be obtained by cleaving \(\beta\)-Bi\(_4\)I\(_4\) samples using tape. e | The left part shows a real-space image of the measured \(\beta\)-Bi\(_4\)I\(_4\) sample taken using an optical microscope. The middle part shows a photoemission intensity map of the same \(\beta\)-Bi\(_4\)I\(_4\) surface, and the right part an enlargement of an area of this image. The white circle (with a diameter \(< 1 \mu m\)) indicates the selected position for nano-ARPES measurements. f | ARPES intensity plots (bottom), and the momentum distribution curves (blue curves, top) at the Fermi level (\(E_F\)) along the \(k_x\) direction at the \(\Gamma\) (left) and \(\bar{Z}\) (right) points of the side (100) surface. The black dotted curves are the Lorentzian fitting of the two peak structures (indicated by red bars) in the momentum distribution curves. \(E\), energy. Panel a is adapted with permission from ref.\(^{39}\), International Union of Crystallography. Panels b-f are adapted from ref.\(^{39}\), Springer Nature Limited.

...recent application of spatially resolved ARPES in the study of \(\beta\)-Bi\(_4\)I\(_4\), the first experimentally verified weak topological insulator.

In 3D, a topological insulator can be classified as either strong or weak depending on the \(Z_2\) topological invariants.\(^{22a,22b}\) A strong topological insulator state, which manifests as gapless topological surface states at all the surfaces, has been experimentally confirmed in many materials, such as Bi\(_n\)Sb\(_m\)\(^{22c}\) and Bi\(_n\)(Se,Te)\(_m\)\(^{43,44,73–75}\). By contrast, the weak topological insulator state is very challenging to detect, because the corresponding gapless surface states emerge only on particular surfaces, which are undetectable in most 3D crystals. Recently, using nano-ARPES, direct experimental evidence was obtained for the existence of a weak topological insulator state in \(\beta\)-Bi\(_4\)I\(_4\) through the observation of the topological Dirac surface states on the side (100) surface.\(^{39}\) \(\beta\)-Bi\(_4\)I\(_4\) exhibits no topological surface states on the (001) surface but exhibits quasi-1D topological surface states on the side (100) surface. These topological surface states exhibit two Dirac cones at the \(\Gamma\) and \(\bar{Z}\) points, which, in principle, can be directly detected by ARPES. However, \(\beta\)-Bi\(_4\)I\(_4\) single crystals are typically very small (\(\sim 30 \mu m\)) along the stacking direction (\(c\) axis), and the cleaved (100) surface is composed of several small domains or stages (typically \(\sim 2 \mu m\) in size). Therefore, it is difficult to study the surface states on the (100) surface by conventional ARPES, as the spot size (\(\sim 100 \mu m\)) of the incident beam is much larger than 2 \(\mu m\). To selectively examine the Dirac surface states, a spatially resolved ARPES measurement with a beam spot \(< 1 \mu m\) was performed on the brightest region of the (100) surface (indicated by the white circle in the right part of Fig. 7f). This approach...
made it possible to resolve the gapless Dirac dispersions expected at the $\Gamma$ and $\bar{Z}$ points [Fig. 7], hence demonstrating the experimental realization of the weak topological insulator state in $\beta$-Bi$_4$I$_3$.

**Conclusion**

With the unique capability to directly visualize and discriminate surface and bulk electronic states, modern ARPES has had a central role in the study of topological materials.

Efforts over the past 30 years have dramatically improved the capabilities of ARPES. Achievements include the introduction of new cryostats, such as $^3$He cryostats [Ref. 1], to enable the study of materials (such as low-$T_c$ superconductors) at sample temperatures <1 K. Moreover, the development of new types of electron analysers, such as DA30-L and ARTOF analysers, has increased performance and data acquisition efficiency, enabling the detection of 2D in-plane electronic structure at a fixed sample geometry. In terms of light sources, many advanced synchrotron-based VUV and soft-X-ray beamlines are now available, including the Dreamline at the Shanghai Synchrotron Radiation Facility (China), which has a high energy resolution (∼25 meV at 1 keV) and a wide photon energy range (20–2,000 eV), and the I05 beamline at Diamond (UK) and beamline 7.0.2 [Ref. 2] of the Advanced Light Source at the Lawrence Berkeley National Laboratory (USA), both of which have nanoscale beam spot sizes. The past 10 years has also witnessed the emergence of a new light source: X-ray free-electron lasers (XFELs) [Ref. 3]. XFELs may be a promising light source for time-resolved ARPES owing to their tunable light spectrum and femtosecond-pulsed beams with a high photon flux. In particular, the wide tunability in energy makes it possible to perform systematic time-resolved ARPES measurements over the full 3D Brillouin zone, which is crucial for investigating the ultrafast dynamics of the bulk electronic structure of 3D materials such as topological semimetals. Nevertheless, the low repetition rate (usually tens to hundreds of hertz) of most XFELs is a major restriction for the application of time-resolved ARPES, because the pulse intensities must also be kept low to minimize space-charge effects. Therefore, the photoelectron count is not yet sufficient for acquiring time-resolved spectra within a reasonable time frame. However, the recently developed European XFEL (Germany) has a high repetition rate (up to 27 kHz) [Ref. 4], and the upgrade to the existing free-electron-laser source at the Linac Coherent Light Source (LCLS-II, USA) to achieve a repetition rate of 1 MHz is expected to commission in 2020. These developments thus provide a promising platform for performing next-generation time-resolved ARPES measurements with high-repetition-rate X-ray photons.

Looking forward, there remains substantial scope to improve the capabilities of ARPES. High-momentum and energy-resolution laser ARPES with $^3$He cryostats is needed to probe the superconducting gap of low-$T_c$ superconductors; this technique is already under development in several laboratories. Despite great interest, spin-resolved soft-X-ray ARPES has still not been realized, owing to the poor efficiency of the spin detectors and the low cross sections of soft-X-ray photoemission. The development of high-efficiency spin detectors and ultrabright synchrotron light will make it possible to perform spin-resolved ARPES measurements in the soft-X-ray regime. Indeed, scientists at several advanced soft-X-ray beamlines, such as the Advanced Resonant Spectroscopy beamline at the Swiss Light Source and the Dreamline at the Shanghai Synchrotron Radiation Facility, have already started working on adding high-efficiency spin spectrometers to their ARPES systems [Ref. 5].

With continued improvements in the focusing optics, we also expect that the efficiency of spatially resolved ARPES will improve. For time-resolved ARPES, it remains necessary to develop high-repetition-rate and stable laboratory-based lasers with tunable probe photon energy and high photon flux. At the same time, the development of strong tunable pump pulses from the ultraviolet down to the terahertz range is essential for the study of coherent and/or resonant excitations of low-energy modes, such as lattice vibrations, and various quasiparticle excitations in non-trivial band structures.

Indeed, the first time-resolved ARPES system with a terahertz pump laser was reported in 2018 [Ref. 6] and successfully applied to the study of Dirac fermions in Bi$_2$Te$_3$. Furthermore, the combination of high-efficiency spin detectors with time-resolved ARPES systems would provide new opportunities for simultaneously acquiring time-resolved and spin-resolved ARPES spectra. For example, a high-efficiency ARTOF–VLEED spin spectrometer that combines a VLEED spin detector with an ARTOF analyser has been developed [Ref. 7]. This high-efficiency spectrometer enables simultaneous spin-resolved, time-resolved and angle-resolved photoemission with pulsed lasers [Ref. 8]. Lastly, combining ARPES with advanced methods of sample synthesis (for example, molecular beam epitaxy and pulsed laser deposition) and other complementary in situ characterization techniques (such as a scanning tunneling microscopy, optical spectroscopy and transport measurements) would provide an efficient platform to systematically design and synthesize new materials, as well as to investigate their physical properties.

Although ARPES has facilitated the discovery and understanding of the unique physical properties of many topological materials, including topological insulators, topological Dirac and Weyl semimetals and unconventional fermions, there are still several prominent physics problems waiting to be explored using this technique. For example, the spin texture of bulk Weyl cones has not yet been investigated owing to the lack of spin-resolved soft-X-ray ARPES systems. Moreover, many predicted topological phases, such as magnetic Weyl semimetals [Ref. 9], still await strong electronic-structure evidence from ARPES to establish their existence. The recent discovery of superconductivity in twisted bilayer graphene [Ref. 10] has attracted extensive attention, and the interesting electronic structure might be resolved by spatially resolved ARPES. Hence, we strongly believe that ARPES will continue to have a leading role in the research into topological materials in the future.

Published online 27 August 2019
83. Young, S. M. et al. Dirac semimetal in three dimensions. Phys. Rev. Lett. 104, 146405 (2010).
84. Wang, Z. et al. Dirac semimetal and topological phase transitions in A\textsubscript{1}Bi\textsubscript{2} (A = Na, K, Rb). Phys. Rev. B 85, 195103 (2012).
85. Wang, Z., Weng, H., Wu, Q., Dai, X. & Fang, Z. Three-dimensional Dirac semimetal and quantum transport in TaAs. Phys. Rev. B 88, 165101 (2013).
86. Liu, Z. et al. A stable three-dimensional Dirac semimetal Ca\textsubscript{3}As\textsubscript{3}. Nat. Mater. 13, 677–681 (2014).
87. Liu, Z. et al. Discovery of a three-dimensional Dirac semimetal Na\textsubscript{4}Bi. Science 343, 866–867 (2014).

Experimental study reporting the discovery of 3D massless Dirac fermions in a topological Dirac semimetal.

88. Soluyanov, A. A. et al. Type-II Weyl semimetal. Science 343, 865–866 (2014).
89. Xu, G., Weng, H., Wang, Z., Dai, X. & Fang, Z. Chern semimetal and the quantized anomalous Hall effect in HgGe\textsubscript{3}. Phys. Rev. Lett. 107, 186806 (2011).
90. Huang, S. M. et al. A Weyl fermion semimetal with surface fermi arcs in the transition metal monopnictide Ta\textsubscript{3}As. Nat. Commun. 6, 7375 (2015).
91. Lu, B. et al. Experimental discovery of Weyl semimetal Ta\textsubscript{3}As. Phys. Rev. X 5, 01105 (2015).
92. A pioneering experimental ARPES study that reports the discovery of Weyl fermion and Fermi arc surface states in the electronic structure of pyrochlore iridates.

93. Liu, Z. K. et al. A stable three-dimensional topological semi-metal and Fermi-arc surface states in Fe\textsubscript{3}As\textsubscript{2}. Nature 495–498 (2015).
94. Young, S. M. et al. Dirac semimetal in three dimensions. Nat. Phys. 8, 116–119 (2012).
95. Lu, D. et al. Angle-resolved photoemission studies of Fe\textsubscript{3}As\textsubscript{2}. Nat. Mater. 12, 945–950 (2013).
96. Yang, H. et al. Visualizing electronic structures of Fe\textsubscript{3}As\textsubscript{2}. Nat. Mater. 12, 945–950 (2013).
97. Vilmercati, P. et al. Evidence for three-dimensional superconducting gap function in optimally doped Fe\textsubscript{3}As\textsubscript{2}. Phys. Rev. B 84, 195148 (2011).
98. Weng, H. et al. High-intensity xenon plasma discharge lamp for laser-based angle-resolved photoelectron spectroscopy. Rev. Sci. Instrum. 87, 113907 (2016).
99. Heinzmann, U. & Dil, J. H. Spin–orbit-induced Fermi arcs and chiral anomalies. Nat. Mater. 15, 1144–1145 (2016).
100. Kirmse, T. & Feder, R. Spin polarization in double diffraction of low-energy electrons from W(001) experiment and theory. Phys. Rev. Lett. 120, 046101 (2018).
101. Lingueg, P., Klein, D. T. & Colotta, R. Low-energy diffuse scattering electron-spin polarization analyzer. Rev. Sci. Instrum. 88, 053106 (2017).
102. Tillmann, D., Thiel, C. & Wolter, E. Very-low-energy spin-polarized electron diffraction from Fe(001). Z. Phys. B 77, 1–2 (1989).
103. Winkelmann, A., Hartmann, C., Engelhard, H., Chiang, T.-C. & Kirmse, J. High efficiency electron spin polarization analyzer based on exchanging scattering at Fe(001). Rev. Sci. Instrum. 79, 83505 (2008).
104. Joziuk, C. et al. A high-efficiency spin-resolved photoemission spectrometer combining time-of-flight spectroscopy with exchange-enhanced spin-polarimetry. Rev. Sci. Instrum. 81, 53904 (2010).
105. Souma, S., Takayama, A., Sugawara, K., Sato, T. & Takahashi, T. Ultrahigh-resolution spin-resolved photoemission spectrometer. Rev. Sci. Instrum. 81, 35104 (2010).
106. Nakada, K. et al. Spin-polarization detection. Phys. Rev. Lett. 107, 207601 (2011).
107. Strocov, V. N., Petrov, V. N. & Dil, J. H. Concept of a multichannel spin-resolving electron analyzer based on Mott scattering. J. Synchrotron Radiat. 23, 708–716 (2015).
108. Joziuk, C. et al. Wider spread spin polarization effects in photoemission from topological insulators. Phys. Rev. B 84, 165112 (2011).
109. Heinzmann, U. & Dil, J. H. Spin–orbit-induced photoelectron spin polarization in angle-resolved photoemission from both ferromagnetic target materials. J. Phys. Condens. Matter 24, 175001 (2012).
110. Osterwalder, J. Can spin-polarized photoemission measure spin properties in condensed matter? J. Phys. Condens. Matter 24, 171001 (2012).
111. Westphal, C., Bansmann, J., Getzlaff, M. & Schönhense, G. Circular dichroism in the angular distribution of photoelectrons from oriented CO molecules. Phys. Rev. Lett. 63, 151–154 (1989).
112. Schneider, C. M. & Kirmse, J. Spin- and angle-resolved photoemission from doped metal surfaces with circularly polarized light. Crit. Rev. Solid State Mater. Sci. 20, 179–285 (1995).
113. Wang, H. et al. Observation of spin-dependent helical spin texture in In\textsubscript{13}S\textsubscript{4}Se\textsubscript{20} from circular dichroism angle-resolved photoemission spectroscopy. Phys. Rev. Lett. 107, 206402 (2011).
114. Weisheit, M. Time-resolved two-photon photoemission from metal surfaces. J. Phys. Condens. Matter 14, R1099–R1141 (2002).
