Flux-flow resistivity anisotropy in the instability regime in the \(a\)-\(b\) plane of epitaxial \(YBa_2Cu_3O_{7-\delta}\) thin films

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Measurements of the nonlinear flux-flow resistivity \(\rho\) and the critical vortex velocity \(v_\phi^c\) at high voltage bias close to the instability regime predicted by Larkin and Ovchinnikov are performed along the node and antinode directions of the d-wave order parameter in the \(a\)-\(b\) plane of epitaxial \(YBa_2Cu_3O_{7-\delta}\) films. In this pinning-free regime, \(\rho\) and \(v_\phi^c\) are found to be anisotropic with values in the node direction larger on average by 10\% than in the antinode direction. The anisotropy of \(\rho\) is almost independent of temperature and field. We attribute the observed results to the anisotropic quasiparticle distribution on the Fermi surface of \(YBa_2Cu_3O_{7-\delta}\).

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![Graph](image_url)

**FIG. 1:** (Color online) A schematic drawing of the bridges on the wafers, and a node and antinode bridges voltage \(V\) versus current \(I\) at \(B = 6\) T and \(T = 73\) K. Inset (a) shows the corresponding \(I\) versus \(V\) just below the maximum of \(I\) versus \(V\) with fits using Eq. (1), inset (b) shows the dynamic nature of \(R_1\) and \(R_2\), see text.

Flux flow due to the Lorentz force in a current carrying type-II superconductor in the mixed state has been studied intensively for many years. In view of the \(d_{x^2-y^2}\)-wave symmetry of the pair wavefunction in the hole-doped cuprates, measurements of the flux-flow resistance (FFR) as a function of the crystallographic direction within the \(CuO_2\) planes become highly interesting. In \(c\)-axis oriented epitaxial \(YBa_2Cu_3O_{7-\delta}\) (YBCO) films under an applied magnetic field parallel to the \(c\)-axis, one could envision an anisotropic behavior resulting in some difference between the flux motion in the node and antinode directions. However, up to now no such anisotropy has been reported, nor has such an effect been analyzed theoretically. Recently, a search study for flux flow resistivity anisotropy in the low bias pinning regime was carried out by some of us in epitaxial thin films of YBCO. To within the experimental error in that study, no anisotropy was found between the transport properties of microbridges patterned along the node and antinode directions. Since at low bias the transport results are determined by the pinning properties of the films, it was concluded that the pinning properties in YBCO are isotropic. At high bias however, close to the Larkin Ovchinnikov (LO) instability regime, pinning is negligible and flux flow resistivity measurements could reveal an intrinsic anisotropy between the node and antinode directions. In this study we present the first experimental results indicating such anisotropy in the \(a\)-\(b\) plane of YBCO films, where the nonlinear FFR along the node direction is enhanced on average by about 10\% compared to that along the antinode direction.

Our measurements were performed on the same two wafers used previously in the low bias measurements. We used two high quality epitaxial, \(c\)-axis oriented YBCO films of 0.12 \(\mu m\) thickness prepared under identical conditions by laser ablation deposition on \((100)\) \(SrTiO_3\) (STO) wafers of \(10 \times 10\ mm^2\) area. The films were patterned with a photolithographic mask into 10 equally spaced \(0.12 \times 12 \times 100\ \mu m^3\) bridges along the line dividing the wafer into two halves by deep UV photolithography and Ar ion milling. The dimensions of the resulting bridges were measured with optical and atomic force microscopes and found to conform with the mask design to within \(\pm 0.4\%\). Successive microbridges were oriented at alternating angles of \(0^\circ\) and \(45^\circ\) to the edge of the wafers as shown schematically in the upper panel of Fig. 1, so that the transport current would flow either along the node or the antinode directions of the order parameter. The alternating direction of adjacent bridges is important in order to minimize systematic differences due to possible large-scale inhomogeneities in the films. On one film, where the orientation of the STO substrate was with the edge of the wafer parallel to the \(c\)-axis, five odd number bridges were along the
antinode direction, and five even number bridges along the node direction. In the other film where the side of the wafer was parallel to the (110) orientation, the role of the antinode and node bridges was reversed due to the epitaxial growth of the film. Studying these two types of wafers was done in order to check if our ion milling process, done at an incident angle of 45° to the wafers, is affecting the properties of the bridges. Any observed difference in the transport properties of the two wafers would imply that the effect is not intrinsic, and results from the patterning process. Low resistance gold contacts were used, and the wafers were annealed for optimal oxygen doping of the films. The resistance measurements were done by the standard four-point dc technique in Helium gas environment with and without a magnetic field of up to 8 T normal to the wafers.

The low bias normal-state resistivity $\rho$ under zero-magnetic-field was measured versus temperature for each bridge on the two wafers. Averaged values of the innermost bridges are given in Figs. 1 and 2 of Ref. 5. The onset of the superconducting transition at $B = 0$ was identical for the two kinds of bridges on the two wafers, with $T_{\text{onset}}^c = 94$ K. This value of $T_{\text{onset}}^c$ has also been reported by Gross et al. for the same kind of films. The resistivity in the normal state at 95 K was $\rho(95$ K$) = 168 \mu\Omega \text{cm}$, identical for both wafers and both type of bridges. The observation that $\rho(\text{node}) = \rho(\text{antinode})$ results from the four-fold symmetry and the heavy twinning of our films.

At lower temperatures, the FFR in the pinning regime of YBCO was recently found to be isotropic. To determine the pinning-independent FFR we measured the current-voltage characteristic near the LO instability caused by the nonequilibrium distribution of the quasiparticles at moderate electric fields. This instability has been studied in detail both in conventional superconductors and in the cuprates. Following procedures of earlier experiments, for each temperature and magnetic field, the voltage applied to the sample was ramped up until a resistive jump due to the LO instability was reached. A typical such jump is seen in the I-V curves of Fig. 1. Above this jump a thermal runaway is observed (not shown here), where the bridges reach the normal state with a typical resistivity of 200 $\mu\Omega \text{cm}$ which corresponds to about $\rho(105$ K$)$. We normally did not expose our bridges to this regime to avoid possible thermal damage. Eq. (53) of the LO theory yields

$$I = \frac{V}{R_{1}} \left\{ 1 + c \sqrt{1 - \frac{T}{T_c}} \right\} - \frac{V}{R_2} \left( \frac{V}{V^*} \right)^2 \left\{ 1 + \left( \frac{V}{V^*} \right)^2 \right\}^{-1}$$

where $c$ is a number of order unity, $V^*$ is a critical voltage, and $R_1$ and $R_2$ are two FFRs which were taken equal in the original theory. We stress that Eq. (1) is valid only in the flux flow regime at high bias near the maximum of I versus V, and not in the pinning regime at low bias. In the low bias regime where $I$ is linear in $V$, the FFR was found to be isotropic, as is expected for a heavily twinned YBCO crystal with a four-fold symmetry if the pinning is also isotropic. By fitting our I-V data in the instability regime using Eq. (1) we find that the FFR $R_1/(1 + \sqrt{1 - T/T_c})$ and $R_2$ have similar values to within 10%. Since the Lorentz force is oriented perpendicular to the current and the applied magnetic field, and YBCO has a four-fold symmetry, flux flow is along the node direction in the node bridges, and in the antinode direction in the antinode bridges. The resistive voltage across the bridge is $V = |\mathbf{v}_F \times \mathbf{B}| L$, where $\mathbf{v}_F$ is the flux-flow velocity, $\mathbf{B}$ is the magnetic field and $L$ is the bridge length. The critical voltage $V^*$ in Eq. (1) is thus associated with a critical vortex velocity $v^*_c$ by this relation.

Joule heating of our samples could affect our results. The possible temperature rise is $\Delta T = \alpha P$, where $P$ is the dissipated electric power per unit area of the sample, and $\alpha$ is the thermal boundary resistance. YBCO films on SrTiO$_3$ substrates have $\alpha \approx 10^{-4}$ cm$^2$ K/W at 300 K. Assuming an even larger value of 10$^{-3}$ cm$^2$ K/W, and taking the maximum values of $P$ close to the LO instability, we estimate that $\Delta T$ remains smaller than a few tenth of a Kelvin. We therefore conclude that Joule heating of our samples can be neglected, except perhaps close to $T_c$. Experimentally, different rates of the current sweeps (14 - 300 mA/s) showed reproducibility of the I-V data, thus excluding the possibility of magnetic hysteresis and thermal heating effects.

Two typical I-V curves for a node and antinode bridges are shown in Fig. 1 at $T = 73$ K and $B = 6$ T. The difference between the two curves is very clear and typical of most of our bridges. With increasing current above the critical current regime, the curve becomes steeper and steeper up to the instability. In the case of voltage-controlled operation as we have, a negative differential resistance sets in just before the jump as seen in Fig. 1. In the pinning-independent instability regime, the flux-flow process follows closely Eq. (1). Indeed, in inset (a) to Fig. 1 where ten data points of $I$ versus $V$ just below the maximum of $I$ are plotted for each curve, one can see that the fits using Eq. (1) are in excellent agreement with the data. In the fits we assumed that $c = 1$, and obtained $R_1$, $R_2$, and $V^*$. In the following we use the procedure as in inset (a) to Fig. 1, to present our results.

Plots of the flux flow resistivities $\rho_1/(1 + \sqrt{1 - T/T_c})$ and $\rho_2$ (calculated from $R_1$ and $R_2$) versus the magnetic field B are shown in Figs. 2 and 3 for the two wafers, respectively, at various temperatures. The spread of the mid-point $T_c$ (where the low bias $dR/dT$ is highest) of the different bridges on the two wafers are given in the insets to these two Figures. The missing data for bridges #4 and #9 in Fig. 2 is due to bad contacts. The relatively large spread of the $T_c$ values on this wafer of about 1.2 K, and the asymmetry relative to the center of the wafers is due to an off-center alignment of the laser ablated plume on the wafer. Nevertheless, bridges #5-8 on LAF551 (Fig. 2) and bridges #3-6 on LAF552 (inset
that mainly to this averaging procedure. These figures show measured node bridges and the two measured antinode directions. The inset shows the distribution of the $T_c$ values on the wafer.

In Figs. 2 and 3 average values are shown of the two $R$ values was found. In this figure $R^2$ is the correlation with the corresponding $T_c$ values for all fields and temperatures.

FIG. 2: (Color online) Averaged flux flow resistivity $\rho_1/(1 + \sqrt{1 - T/T_c})$ versus magnetic field $B$ on LAF551 at 60, 70 and 78 K, for the node (stars) and antinode (squares) directions. The inset shows the distribution of the $T_c$ values on the wafer.

FIG. 3: (Color online) Averaged flux flow resistivity $\rho_2$ versus magnetic field $B$ on LAF552 at 67, 73 and 77 K, for the node (stars) and antinode (squares) directions. The inset shows the distribution of the $T_c$ values for the four measured bridges on the wafer.

of Fig. 3, in a reversed order of the bridges) have a much smaller $T_c$ spread of about 0.4K. Thus the FFR in these bridges was measured and compared, and no systematic correlation with the corresponding $T_c$ values was found. In Figs. 2 and 3 average values are shown of the two measured node bridges and the two measured antinode bridges on each wafer. The errors in these figures are due mainly to this averaging procedure. These figures show that $\rho_1/(1 + \sqrt{1 - T/T_c})$ and $\rho_2$ increase almost linearly with $B$ at low temperatures, but more than linearly at higher temperatures. In all cases, a clear anisotropy is observed where the $\rho_1/(1 + \sqrt{1 - T/T_c})$ and $\rho_2$ values along the node direction are larger than along the antinode direction. The fact that the linear term in Eq. (1) $R_1 \propto \rho_1$ was found to be anisotropic seems to contradict the four-fold symmetry of the twinned crystal. We note, however, that $R_1$ (and also $R_2$ and $V^*$) are voltage dependent (see inset (b) to Fig. 1). The results of this inset were obtained by fitting the antinode $I$-$V$ data using Eq. (1) near the instability regime at various voltage segments of $V \pm 5$ mV. The segments contained between 6 data points (above the maximum of $I$ and just below the voltage jump) to 17 points at the lowest voltage bias shown. All fits had $R^2$ of better than 0.999 (solid symbols). Also shown, are results of two inferior fits with 48 and 65 data points (empty symbols), and as one can see, the resulting parameters still follow the general behavior. We thus conclude, that $R_1$ is not a linear term of the FFR, and therefore does not have to obey the symmetry of the heavily twinned crystal.

Next, we turn to the critical voltage $V^*$ which corresponds to the critical vortex velocity $v^*_\phi$ according
to \( V^* = |v^*_v \times B|/L \). In Fig. 5 we plot average values of the critical vortex velocity \( v^*_v \) of the node (6,8) and antinode (5,7) bridges on LAF551 versus magnetic field. In the node direction \( v^*_v \) is systematically larger (by about 10-20\%) than in the antinode direction. In LAF552 this anisotropy is similar (not shown). The \( v^*_v \) values at each temperature show only a weak field dependence, and are basically quite constant as was found previously. Our conclusion therefore is that the anisotropy in \( v^*_v \) is consistent with the anisotropy in \( \rho_1 \) and \( \rho_2 \), as all have comparable percentage values.

In the absence of a theoretical prediction for the anisotropies of the flux flow resistivities \( \Delta \rho_1 \) and \( \Delta \rho_2 \), and the vortex velocity \( \Delta v^*_v \), we shall present simple qualitative arguments for explaining them. In principle, in the absence of pinning, the vortex velocity \( v^*_v \) is affected by the ability of the vortices to transfer momentum to the quasiparticles. The Magnus or Lorentz force that acts on a vortex constantly transfers momentum to it. After a short acceleration, the vortex reaches a terminal velocity \( v^*_v \) in the viscous media while the extra momentum has to be constantly dissipated to the quasiparticle excitations. Thus the easier it is for a vortex to transfer momentum when it moves in a certain direction, the faster it will move in this direction. It was predicted theoretically that in a d-wave superconductor the low energy excitations are located in the vicinity of the nodes. This was verified in photoemission experiments where four small regions of quasiparticle excitations were found on the Fermi surface in the directions of the nodes. Our experimental observations indicate larger flux flow velocity and FFR in the node direction. Therefore, it seems that the momentum transfer by vortices to the quasiparticles is more effective in the node direction as compared to that along the antinode direction where less quasiparticles are available for the momentum transfer process. Clearly, a comprehensive theoretical analysis of this issue is needed in order to make a detailed quantitative comparison with the present results.

Finally, we rule out the possibility that twinning is responsible for the observed anisotropy. Twins in thin YBCO films form a dense mosaic of elongated crystallites with boundaries along the (110) and (1\overline{1}0) node directions as shown schematically in the inset to Fig. 5. If the twin boundaries serve as easy channels for flux flow, then flow in the anti-node bridge is unhindered since there is always a Lorentz force component in the direction of the motion. In the node bridge however, the flow is blocked each time it reaches a boundary normal to its direction. This should lead to a larger flux flow velocity and FFR in the antinode direction as compared to the node direction. Since this is opposite to the observed results, twins could not be the source of our observations.

In summary, we observed a clear anisotropy in the nonlinear flux flow resistivities \( \rho_1 \) and \( \rho_2 \) of YBCO which we attribute to the intrinsic anisotropy of the low energy quasiparticle distribution on the Fermi surface of a d-wave superconductor. The observed anisotropy in the critical vortex velocity \( v^*_v \) is consistent with the anisotropy of the measured flux-flow resistances.

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1. A. I. Larkin, and Yu. N. Ovchinnikov, Z. Eksp. Teor. Fiz. 68, 1915 (1975) [Sov. Phys. JETP 41, 960 (1976)].

2. R. P. Huebener, Magnetic Flux Structures in Superconductors, (Springer, Berlin, 2nd edition, 2001).

3. M. Tinkham, Introduction to Superconductivity (McGraw-Hill, 2nd edition, 2001).

4. C. C. Tsuei and J. R. Kirtley, Rev. Mod. Phys. 72, 969 (2000).

5. G. Koren, P. Aronov and E. Polturak, Phys. Rev. B 72, 212503 (2005).

6. R. Gross, P. Chaudhari, D. Dimos, A. Gupta, and G. Koren, Phys. Rev. Lett. 64, 228 (1990).

7. S. G. Doettinger, R. P. Huebener, R. Gerdemann, A. Kuehle, S. Anders, T. G. Traueble, and J. C. Villegier, Phys. Rev. Lett. 73, 1691 (1994).

8. C. D. Marshall et al., J. Appl. Phys. 73, 850 (1993).

9. B. Kalisky, Y. Wolfus, Y. Yeshurun, G. Koren, and R. P.
Huebener, Physica C 401, 273 (2004).

10 G. E. Volovik, Pis’ma Zh. Eksp. Teor. Fiz. 58, 457 (1993), [JETP Lett. 58, 469 (1993)].

11 M. R. Norman et al., Nature, 392, 157 (1998).