SOURCE-FREQUENCY PHASE-REFERENCING OBSERVATION OF AGNs WITH KAVA USING SIMULTANEOUS DUAL-FREQUENCY RECEIVING

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Abstract: The KVN(Korean VLBI Network)-style simultaneous multi-frequency receiving mode is demonstrated to be promising for mm-VLBI observations. Recently, other Very long baseline interferometry (VLBI) facilities all over the globe start to implement compatible optics systems. Simultaneous dual/multi-frequency VLBI observations at mm wavelengths with international baselines are thus possible. In this paper, we present the results from the first successful simultaneous 22/43 GHz dual-frequency observation with KAVA(KVN and VERA array), including images and astrometric results. Our analysis shows that the newly implemented simultaneous receiving system has brought a significant extension of the coherence time of the 43 GHz visibility phases along the international baselines. The astrometric results obtained with KAVA are consistent with those obtained with the independent analysis of the KVN data. Our results thus confirm the good performance of the simultaneous receiving systems for the non-KVN stations. Future simultaneous observations with more global stations bring even higher sensitivity and micro-arcsecond level astrometric measurements of the targets.

Key words: galaxies: active — galaxies: individual: 4C 39.25, 0945+408 — astrometry — techniques: interferometric

1. INTRODUCTION

Very long baseline interferometry (VLBI) observations at millimeter wavelengths are conventionally difficult due to tight limitations on the integration time imposed by fast phase fluctuations caused by atmospheric propagation effects. The unique quasi-optical (hereafter QO) system implemented on the Korean VLBI Network (KVN) that enables simultaneous multi-frequency receiving overcomes some of these difficulties (Han et al. 2008, 2013; Jung et al. 2015). The non-dispersive characteristics of tropospheric propagation effects (Rioja & Dodson 2011), which dominate the phase noise at high frequencies, lead to strong correlations between visibility phases at different frequencies (e.g., Jung et al. 2011). It is thus possible to calibrate these effects in the high frequency phases with their lower frequency counterparts (e.g., Middelberg et al. 2005; Rioja & Dodson 2011; Jiang et al. 2018). Simultaneous multi-frequency receiving has made this kind of calibration very feasible as no temporal interpolations are necessary (e.g.,...
Rioja & Dodson 2011; Rioja et al. 2015; Algaba et al. 2015). The immediate benefit of applying this method (so-called frequency phase transfer, FPT) is the extension of the coherence time at high frequency. Rioja et al. (2015) found that the coherence time at 130 GHz increased from several seconds to 20 minutes after FPT. It is possible to further extend it to hours if more than 3 frequencies are simultaneously observed (Zhao et al. 2018), making the KVN very powerful in detecting weak sources at high frequencies (e.g., Algaba et al. 2015).

Another benefit of the QO system is improved mm-wave astrometry. The long coherence time and weaker dependence on line of sights of the phases after FPT allow switching between sources in order to calibrate residual propagation effects, i.e., source-frequency phase-referencing (SFPR, Rioja & Dodson 2011). This will also provide astrometric information which is a combination of the relative positions between two frequencies in the two sources (e.g., Rioja & Dodson 2011; Rioja et al. 2014, 2015; Jiang et al. 2018). For an AGN jet, the frequency-dependent shift in the position of the VLBI core is known as the core-shift (Blandford & Königl 1979). KVN is doing regular astrometric observations with SFPR up to 130 GHz (e.g., Rioja et al. 2014, 2015; Dodson et al. 2014, 2018; Yoon et al. 2018), and consistent results were found between KVN and the Very Long Baseline Array (VLBA) (Rioja et al. 2014).

The disadvantages of KVN, however, come mainly from the short baseline lengths and the limited number of baselines. Poor image qualities (e.g., Zhao et al. 2018) and structural blending effects in astrometric measurements (Rioja et al. 2014; Zhao et al. 2015) are often seen in the results. Nevertheless, as the simultaneous observing mode has been demonstrated to have great potential for future mm-VLBI observations (e.g., Dodson et al. 2017), more stations all over the world begin to implement compatible optical systems. Non-KVN stations to have been upgraded with QO systems include the 4 VERA (VLBI Exploration of Radio Astrometry) stations in Japan and the Yebes 40m telescope in Spain. These stations are implemented with KVN compatible dual-frequency QO systems. The data obtained in the second session are yet to be analyzed and the scientific results will be presented in forthcoming papers. Here we summarize the observation and data analysis for 4C 39.25 campaign in the first session.

2. Observation and Data Analysis

Simultaneous 22/43 GHz VLBI observations with the joint array of KVN and the international partner stations were successfully carried out in two sessions, on November 7–9, 2016 and March 16–18, 2018, respectively. The participating antennas included all three KVN telescopes in Korea, the four VERA telescopes in Japan, and the Yebes 40m telescope in Spain. These stations are implemented with KVN compatible dual-frequency QO systems. The data obtained in the second session are yet to be analyzed and the scientific results will be presented in forthcoming papers. Here we summarize the observation and data analysis for 4C 39.25 campaign in the first session.

2.1. Observation

The source 4C 39.25 (RA $= 9h27m03.013938s$, Dec $= +39^\circ 02^\prime 20.851770^\prime\prime$, Epoch $= J2000$) was observed on November 9, 2016 for 6 hours while the target was setting as seen from East Asian stations. During the first session, fringe-detections with Yebes were not successful and only two VERA stations (Mizusawa and Iriki) were implemented with QO systems.

The observing frequencies were 22.1 and 42.8 GHz. The total recording bandwidth was 256 MHz (128 MHz at each band). The signal was further split into 32 MHz IFs. The other two VERA stations (Ogasawara and Ishigaki) were observing at 42.8 GHz only and the data were not used in this analysis. Usually, for dual-frequency calibrations, an integer ratio between frequencies is favored because it is ideal for addressing the $2\pi n_p$ phase-wrapping ambiguities (e.g., Middelberg et al. 2005). However, in practice, this requirement cannot always be satisfied, as for example, the rest frequencies of SiO and water maser transitions have a non-integer ratio. In the case of our observation, the frequency windows are limited by the common frequency range among different stations. Nevertheless, FPT with non-integer ratios can be implemented (e.g., Dodson et al. 2014).

Most of the observing time was for tracking 4C 39.25 with the aim of checking the system stability over a long time and a wide range of elevations. 4C 39.25 is suitable for this purpose as it passed near the zenith of most of the stations during the observation campaign. Short (3 minutes) scans on another source, 0945+408 (RA $= 9h45m40.55s$, Dec $= +40^\circ 39^\prime 44.586920^\prime\prime$, Epoch $= J2000$), were inserted every 20 minutes to obtain astrometric measurements using SFPR. The separation between the two sources on the sky plane is $\sim 5^\circ$.

Typical weather conditions were found during the observation at most of the sites, with sky opacity less than 0.15 and opacity corrected system temperatures
Figure 1. Self-calibrated images of 4C 39.25 (left) and 0945+408 (right) at 22 (upper) and 43 GHz (lower). The beams are shown in the lower left corner. The contour levels start from 3 times the r.m.s level and increase by a factor of 2.
(T_{sys}^{*}) ranging between 100 and 600 K during most of the observing time, except around UT 0 at Tamna and during the last two hours at Iriki station due to bad weather conditions (i.e., sky opacity reaches nearly 0.3 and T_{sys}^{*}>1000 K).

After the observations, data recorded at each site were sent to the Korea-Japan joint correlation center (KJCC) and correlated by the hardware correlator (Lee et al. 2015a). The same correlator is also used for most of the KaVA and EAVN (East Asian VLBI Network) observations.

2.2. Data Reduction

The data analysis for the simultaneous 22/43 GHz campaign was carried out with AIPS and Difmap. First, amplitude calibration was performed in two steps, normalizing the auto-correlation spectra at each correlation interval with AIPS task ACCOR and multiplying the gain curves and the atmospheric opacity corrected system temperatures (T_{sys}^{*}) with task APCAL. A correction factor of 1.33 for digital sampling loss and quantization loss was also applied with APCAL (Lee et al. 2015b).

Phase calibration was conducted in three parallel passes. First, we performed global fringe-fitting, self-calibration and hybrid mapping for each source at each frequency independently following the standard procedure in regular imaging observations. The solution intervals for fringe-fitting were set to be the typical coherence times at each frequency, i.e., 60 and 30 seconds for 22 and 43 GHz respectively; Second, we applied SFPR analysis with the data on all baselines of the five stations. We performed fringe-fitting at the lower frequency band, 22 GHz, with a solution interval that matches the typical coherence time of the higher frequency (30s) and derived the phase solutions for the higher frequency band data that are obtained at the same time. We first scaled up the phase solution for the first IF (reference frequency, 22.1 GHz) with a factor of 2 and then calculated the phase solutions for each IF (φ_{IF}) with the multi-band delay (τ_{MB}) and the frequency difference (i.e., φ_{IF} = 2 × φ_{22.1} + τ_{MB} × (ν_{IF} − 2 × 22.1), where φ_{22.1} and ν_{IF} stands for the phase solutions at 22.1 GHz and the frequency of each IF, respectively). After applying the phase solutions derived from the lower frequency (FPT), we estimated the coherence time at the higher frequency band using the simple approach demonstrated in Rioja et al. (2012). We then ran a second fringe-fitting on the 4C 39.25 data at the higher frequency and applied the solutions to both sources by cubic temporal interpolation. Finally, we imaged the 0945+408 data and measured the position offset from the phase center which corresponds to the combination of the relative position between two frequencies of the two sources, i.e., the core-shift; In the third pass, we followed the same procedure as in the second one but with KVN data alone to provide a consistency-check on the reliability of the results.

Figure 2. visibility phases of 4C 39.25 along MIZ baselines at 43 GHz before (dots) and after (crosses) applying FPT from 22 GHz; Each point on the plot corresponds to a temporal average of 10 seconds.
In this section, we present the results of 4C 39.25 & 0945+408 observations during the first international simultaneous 22/43 GHz dual-frequency campaign. The main results include the amplitude self-calibrated dual-frequency maps of both sources, the extension of the coherence time at the higher frequency after applying FPT, and the astrometric measurements with SFPR which is the combined core-shift (or centroid-shift) of the two sources. Each of these results are followed by a brief comparison with either previous KaVA imaging results or the KVN FPT and astrometric results.

### 3.1. Hybrid Maps

KaVA is now conducting imaging observations at 22 and 43 GHz regularly (not simultaneously) for over 1000 hours per year with several large programs running (e.g., Kino et al. 2015). For projects that require dual frequency images, observations at different frequencies are usually done on two successive days. The imaging capability of KaVA was evaluated in details by Niinuma et al. (2014). In this section, we compare the morphology with previous observations by KaVA to check the reliability of the system. Detailed studies of source structure and spectra will be presented together with the result from the second session and follow-up observations in forthcoming papers.

The self-calibrated hybrid maps of the two sources at two frequencies are shown in Figure 1. Our observation reveals similar structures of the targets as compared with previous works. For 4C 39.25, our maps show the peculiar core-jet structure of this source, with a bright, complex, optically-thin feature and a relatively dimmer, compact core component at a distance of ~ 3 mas in the west from the brightest feature. This morphology is consistent with previous studies (e.g., Niinuma et al. 2014). However, the resolution in our 22 GHz map is not sufficient to fully resolve the structure of this source as can be seen from the map. For 0945+408, we find a one-sided core-jet structure with the jet extending in the southeast direction. A slight change of the jet axis direction is shown at ~ 2 mas from the core. This morphology is similar as presented elsewhere, for instance in the 15 GHz results by MOJAVE project (Lister et al. 2013). The weak, extended emissions at ~ 7 mas and beyond from the core revealed by MOJAVE are not obvious on our maps, suggesting that these diffuse emissions have steep spectra.

### 3.2. Extension of Coherence Times

One of the most important improvements to come from implementing the QO system is the extension of the coherence time at higher frequencies by applying FPT (e.g., Rioja et al. 2015; Algaba et al. 2015; Dodson et al. 2017). In the second pass of our data analysis, we calibrate the higher frequency data with the fringe solutions obtained at the lower frequency. The 43 GHz visibility phases of 4C 39.25 at Mizusawa baselines before and after this calibration are shown in Figure 2. This figure clearly demonstrates the power of the simultaneous dual-/multi-frequency receiving. As we can see from the plots, before applying FPT, which is equivalent to a single frequency observation, the phases are dominated by propagation errors. The phase rotation can be much larger than 1 radian within several minutes. After applying FPT, the phases are significantly better aligned. The time taken for the phase to rotate around 1 radian approaches several tens of minutes. Such an improvement in the coherence is comparable with those seen in the KVN-only experiments (e.g., Rioja et al. 2014; Jung et al. 2015).

In order to quantify the extension of coherence time, we followed a simple procedure presented in Rioja et al. (2012). We loaded the raw and FPTed data into Difmap and performed self-calibration over a series of integration times. We then imaged the calibrated data and compared the peak flux densities on the map with the one produced with the shortest solution interval, which is 0.1 minutes, much shorter than the typical coherence time at this frequency. In figure 3 we show the percentage of the recovered peak flux density compared with the latter as a function of the integration time for each dataset. The coherence time is defined as the integration time with which the flux recovery level is 0.6 (Rioja et al. 2012, 2015). We found the coherence time at 43 GHz is 1.5 and 60 minutes before and after applying FPT, respectively.

The limitation of coherence time after FPT comes mainly from ionospheric and instrumental propagation effects. These effects are usually slow varying compared with their tropospheric counterparts as can be seen from the figures. With KVN 4-frequency observations, Rioja et al. (2015) reported the coherence time was extended to 20 minutes up to 130 GHz after FPT. At 43 GHz, this value is expected to be even longer because of the smaller frequency ratio and more stable instrumental performance. Our KaVA results are quite consistent with such an expectation for KVN which con-
firms the performance of the QO system for the VERA stations.

We note this is only a simple estimation of the overall coherence time of the observations. A more detailed analysis on the phase coherence at each baseline using Allan standard deviation and coherence function will be presented in a companion paper (Jung et al. 2019).

3.3. Astrometric Measurements

In this section, we present the astrometric results by applying the calibrations of 4C 39.25 FPTed data to those of 0945+408, which correspond to the second pass of our data analysis. An example of the SFPRed phases of 0945+408 along all baselines are shown in Figure 4 and a corresponding map of the source is shown in Figure 5. Point source models for fringe-fitting were used in this case. Compared with the self-calibrated map in Figure 1, although the noise level in Figure 5 is higher, astrometric information of the target is preserved in this map as an offset between the peak position and the map center can be seen in the figure. This offset corresponds to the combined core-shift or centroid-shift of the two sources.

The reliability of KVN astrometric performance was confirmed in Rioja et al. (2014) by comparing with the VLBA results. We compare our results with those obtained from independent parallel analysis on the KVN data only which corresponds to the third pass of our data analysis. We perform our comparison in three different cases with different structure models used for fringe-fitting, 1) with point source models, 2) with the structural models from the imaging analysis on the data, i.e., those in Section 3.1, and 3) with the models from Niinuma et al. (2014) for 4C 39.25 which were obtained with the full KaVA array but the same as case 2) for 0945+408. The three pairs of results are plotted in Figure 6. We note that although the shift is relatively large, which indicates the presence of structural blending effects, the measurements of KaVA and KVN agree well with each other, 2 σ for point source model and less than 1 σ when structure models are used for fringe-fitting. The results thus confirm the reliability of the performance of QO systems for VERA stations. Furthermore, because KaVA has higher angular resolutions and image sensitivities, the accuracy of the measurements (beamsize divided by signal-to-noise ratio) is much higher than those of KVN. For instance, in Figure 5, the major axis size of the synthesized beam is 1.3 mas and the dynamic range is about 47, which result in an astrometric accuracy of 28 µas, while for the case with KVN alone, the corresponding values are 3.1 mas, 32, and 96 µas, respectively.

Structural blending effects are often seen in KVN astrometric results (e.g., Rioja et al. 2014). Although our KaVA observation has higher angular resolution, we still find a large position offset on the SFPRed map. This offset is ~ 1/4 of the 22 GHz beamsize and along the same direction of the major axis of the beam when point source models are used. When the structural models are used for the SFPR analysis, the amount of offset gets smaller (Figure 6). The most likely reason for this large structure blending effect is the low resolution at 22 GHz as described in Section 3.1. Furthermore, as Iriki station is located close to the Korean peninsula, in the current array configuration, Mizusawa is the only station that has long baselines (> 1000 km). This means in each triangle of the closure phases, there is at least one short baseline, which is comparable in length to a KVN baseline. This explains the comparable level of structural blending effect because the standard VLBI data analysis utilize the closure phase during
Simultaneous Dual-Frequency Observations with KaVA

Center at RA 09 48 55.33815100 DEC 40 39 44.5869200
Cont peak flux = 2.4232E-01 JY/BEAM
Levs = 5.100E-03 * (-3, 3, 4.240, 6, 8.490, 12, 16.97, 24, 33.94)

Figure 5. KaVA SFPRed map of 0945+408 at 43 GHz. The beam is 1.28 × 0.70 mas with a position angle of −71°. The grid serves as a visual guide for the offset of the peak of brightness from the center of the map.

global fringe-fitting and self-calibration.

The limitations of the array with the fact that all the long baselines are associated with one single station will be overcome when simultaneous dual-frequency observations are available with the full KaVA array. The participation of Ogasawara and Ishigaki stations will not only double the longest baseline length, but also will provide closure triangles with all baselines longer than 1200 km (e.g., Mizusawa-Ishigaki-Yonsei triangle). From Figure 6, we can notice that when models obtained with the full KaVA array for 4C 39.25 (Niinuma et al. 2014) are used, the offset becomes much smaller than without. If we were to use full-array models for both sources, the offset should become more reasonable. The second observing session performed in March 2018 and future upcoming observations will be able to provide valuable scientific results.

4. SUMMARY AND FUTURE ASPECTS

We have carried out simultaneous 22/43 GHz dual-frequency observations with the KaVA array. Our study has confirmed the performance of the quasi-optics systems implemented on VERA stations, as the results:

- revealed similar structures of the target sources compared with previous single frequency imaging observations.
- confirmed the significant extension of the coherence time at the higher frequency from ~1 minute to several tens of minutes by applying FPT.

Future observations with the full KaVA array and more stations (e.g., in Spain and Australia) will further improve the resolution and sensitivity of the images as well as more accurate astrometric measurements. The dual-frequency simultaneous observing mode will save approximately half of the time for the dual frequency observations. Running large programs with simultaneous observations will bring new scientific outcomes (e.g., astrometry) without losing current objectives with more machine time becoming available for new scientific projects, especially as more stations join the array.

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