THE SCINTILLATION MODELS FOR SIGNAL PROPAGATION THROUGH SATellite IONOSPHERIC CHANNELS

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Abstract. The results of analysis for signals propagating through the ionosphere satellite communication channels with temporal and spatial electron density irregularities in the ionosphere plasma are presented in the article. These electron density irregularities most commonly occur in low-latitude, auroral, and polar regions and refer to the random signal amplitude and phase fluctuations. Occurrence of scintillation is difficult to predict and model due to the variability of its numerous influencing factors, which include solar activities, inter-planetary magnetic field activities, local electric field and conductivity, convection processes, wave interactions. Satellite ionospheric radio waves in P-, L- frequency bands are vulnerable to scintillations that can severely impact the acquisition and tracking process in receivers, causing a degradation in navigation and in information systems (as example, Kospas-Sarsat system) solution accuracy, integrity, and continuity. The widely-used indice to measure ionospheric scintillation activity is the scintillation index for amplitude scintillation (the standard deviation of the received signal power normalized to the average signal power). The values of scintillation index (0.27…0.49) for Cospas-Sarsat channel (406.0-406.1 MHz) are evaluated in the article. The empirical model for probability density for signal amplitude as m-law Nakagami based on these values of scintillation index is used to evaluate error-performance degradation concerning to propagation through free space. The computer simulations for evaluation of these error-performance degradations are performed - the degradation of signal/noise is about 6.8 dB for bit-error 0.001.

Keywords: satellite ionosphere channels, signals, signal scintillations, error-performances, signals, phase shift keying signals, Kospas-Sarsat

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to calculate the statistical characteristics of signal power variations (for example, average power, statistical power moments) and, as a result, estimate the probability characteristics when receiving signals and compare them with the probability characteristics of receiving signals when propagating in free space [10, 11]. Based on this comparison, the corresponding energy losses are estimated, which must be taken into account in the energy budgets of radio lines.

The results of experimental and theoretical studies using these models show the dependence of energy losses due to scintillation on the central frequency, on the spatial distribution of transmitting and receiving devices, solar activity, daily time [3, 4]. Radio lines of the L-frequency range, actively used by satellite systems of global navigation (Glonass, GPS, Galileo, etc.) [12-17], have been studied in sufficient detail. The problem of the generalization, development and addition of these results for the p-frequency radio links, which are also actively used by satellite systems for transmitting information, is relevant. An example is the international satellite search and rescue system Cospas-Sarsat, operating in the frequency range 406.0 ... 406.1 MHz [18].

2. PROBLEM STATEMENT

In Fig. 1 is a diagram explaining the propagation of \( s(t) \) signals over a satellite ionospheric radio link. Studies show that the main contribution to the scintillation of signals is determined by the inhomogeneities of the electron density of the ionospheric layer \( F \), having dimensions \( l \), comparable to the dimensions of the first Fresnel zone (\( \lambda \) is the wavelength of the signals) [1, 4, 15]. For the location of the onboard transmitter at an altitude of \( H_0 >> H \) (\( H = 350-400 \) km - the height of the ionospheric layer \( F \)), which is a valid condition for satellite global navigation systems (\( H_0 \approx 19,200 \) km), the incident signals \( s(t) \) refer to the far radiation zone and rely flat waves. In models [1, 4], the effect of ionospheric inhomogeneities is given by the action of a thin screen with a field of random phase distribution. The amplitude \( A \) of the resulting signal \( s'(t) \) at the input of the ground receiving point is a random variable and is determined by applying diffraction methods taking into account a random field [1, 3].

For the basic model of channel noise in the form of additive white Gaussian noise \( n(t) \) (ABGN), the rule of optimal reception of digital signals that implements the statistical maximum likelihood criterion is based on the calculation of the cross-correlation of the input implementation \( s'(t) + n(t) \) with the original signal \( s(t) \) [10]. In this case, the probability of erroneous reception of the \( Pb \) bit in the coherent reception of signals with two-phase and four-phase keying
(FM2 signals, FM4 signals) without error-correcting coding is determined by the relation [10]
\[ P_b(E_b / N_0) = 1 - F(\sqrt{2E_b / N_0}). \] (1)

Here, \( E_b \) is the signal energy per information bit; \( N_0 \) is the ABGSh spectral density (one-sided);
\[ F(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \exp(-t^2 / 2)dt. \]

Assuming the amplitude \( A \) of the signal \( s'(t) \) is random stationary in a broad sense of magnitude with the distribution density \( p(A) \), the average error probability taking into account scintillation can be calculated using the expression [10]
\[ P_b = \int_0^\infty P_b(E_b, A^2 / N_0)p(A)dA. \] (2)

The essence of the task is to provide descriptions of the scintillation models of the amplitudes of signals as a random process during propagation through heterogeneous satellite ionospheric radio links, to give the results of calculating the error probability for signals with phase shift keying taking into account scintillation models and to give these signals for the P-frequency range of relation to distribution in free space.

3. MODELS OF SIGNAL SCINTILLATION

The considered scintillations of signals and their statistical characteristics are determined by a number of parameters — the center frequency \( f_0 \), the spatial distribution of the transmitting and receiving devices, solar activity, the speed of the ionospheric irregularities, the daily time, etc. [3, 4].

When creating and developing signals scintillation models taking into account these factors, two approaches are used - based on the use of analytical methods for describing the propagation of signals with ionospheric irregularities and on the use of empirical relationships regarding the distribution density \( p(A) \) [1, 5].

In the approach based on analytical methods for describing the propagation of signals, the presence of ionospheric irregularities with small spatial variations in electron density with a scale comparable to the wavelength \( \lambda \) [1, 4] is considered. Its temporal variations with a \( 1 / f_0 \) scale are also considered to be small. In this case, the component of the electric field \( E \), falling vertically on the ionospheric layer (x-axis) during propagation, is a solution of the wave equation [3, 19, 20]
\[ \frac{d^2E(x,t)}{dx^2} + k^2\varepsilon(x,t)E(x,t) = 0. \] (3)

Here \( k = 2\pi / \lambda \) is the wave number for free space, \( \varepsilon(x,t) \) is the dielectric constant of the medium.

Further, it is assumed that the main variations in the dielectric constant \( \delta\varepsilon(x,t) \approx 4\pi reNe / k^2 \) occur in a limited volume with linear dimensions \( l \) (Fig. 1), comparable to the dimensions of the first Fresnel zone \( r_0 \) (re is the electron radius). Under this condition, equation (3) is transformed into a parabolic equation [1, 3]
\[ -j2k \frac{\partial U}{\partial x} + \nabla_i^2U + k^2\varepsilon(x,t)U = 0. \] (4)

Here
\[ u = U \exp(-j\kappa x), \quad \nabla_i^2 = \frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial z^2}. \]

Equation (4) is stochastic, its solution determines the relationship between the random variables \( U \) and \( \varepsilon(x,t) \) and establishes their statistical characteristics (for example, statistical moments, including the fourth moment of the amplitude of the signal \( s'(t) \), used to estimate the range of fading signals).
Equation (4) is stochastic, its solution determines the relationship between the random variables U and ε (x, t) and establishes their statistical characteristics (for example, statistical moments, including the fourth moment of the amplitude of the signal s′(t), used to estimate the range of fading signals).

Equation (4) is non-linear, the problem of finding its solution in the form of closed analytical expressions remains open. Its approximate solutions of Born and Rytov are known as the sum of terms depending on the small parameter ε′(x, t) << 1, where ε(x, t) = 1 + ε′(x, t) [1, 6, 7]. For the approximation ε′(x, t) = 0, the solution (4) in the form of U0(x, t) determines the propagation of signals in free space. The use of the term depending on ε′(x, t) gives the Born approximation U′(x, t) (Debye-Born scattering) corresponding to single scattering under the condition U′/U0 << 1 [7]. For U′/U0 ≈ 1, solution (4) is equivalent to solving the problem with multiple scattering [7].

The Rytov approximation (small perturbation method) is based on the use of the relation ψ = Ln(A), the function is a solution of the equation [1, 6, 7].

\[-2k \frac{\partial \psi}{\partial x} + \nabla^2 \psi + (\nabla \cdot \psi)^2 + k^2 \varepsilon = 0. \tag{5}\]

The solution (5) in the Rytov’s approximation determines the linear relationship between the random functions and and the relationship of the statistical moments of the function Ln(A) to the amplitude of the signal s′(t) and the fluctuations of the electron density. At the same time, it is noted that the considered analytical approaches using approximate solutions of equation (4) do not provide sufficiently accurate results in the general case of the U′/U0 ratio with respect to experimental measurements [1].

More accurate scintillation models from the second class associate the parameters of empirical distribution densities p(A) of the amplitude of the signal s′(t) with the scintillation index S4 = (<I^2> - <I>^2) / <I>^2, the values of which define an important characteristic for applications - the fading range of Pf signals [5, 9]. Here I = A^2(t) is the signal power; <> Is the averaging operation over an ensemble of signals, or over time, assuming that the random process A(t) is ergodic.

The fading range of the amplitude of the Pf signals (dB) at the input of the receiving device is determined by the approximate relationship [5]. A more accurate relationship between Pf and S4 can be determined using empirical models of the distribution density p(A), the parameters of which can be set using the analytical models considered above (4), (5) or using experimental studies of radio lines.

According to the values of the S4 index, fading is classified: weak to S4 <0.3; average 0.3 <S4 <0.6; strong S4 > 0.6 [5]. An increase in S4 values is accompanied by an increase in energy loss with respect to propagation in free space.

For scintillations, the density p(A) can be represented by well-known distribution laws: the log-normal distribution [2], the m-distribution by Nakagami [5],
the Rayleigh-Rice distribution, the -$\alpha$-$\mu$

distribution [9].

Most often, the amplitude density $p(A)$
of signals during propagation along the
ionospheric radio link is described by the
Nakagami distribution [4, 5]

$$p(A) = \frac{2}{\Gamma(m)} \left( \frac{m}{\sigma^2} \right)^m A^{2m-1} \exp\left( -\frac{mA^2}{\sigma^2} \right).$$ (6)

Here $\sigma^2$ is the fluctuation component of
the signal power $s'(t)$; $m \geq 1/2$ is a parameter
specified by the relation $\Omega = \langle A^2 \rangle$ [21].
The parameters $S_4$ and $m$ are related by the
relation $m = 1 / S_4$ [5].

The distribution by Nakagami is
approximated by the Rayleigh-Rice
distribution [21]

$$p(A) = \frac{A}{\sigma^2} \exp\left( -\frac{A^2 + A_0^2}{2\sigma^2} \right) I_0 \left( \frac{AA_0}{\sigma^2} \right).$$ (7)

Here, $A_0$ is the average amplitude of the
signal component; $I_0(x)$ is a modified Bessel
function of the first kind of zero order [21].

The Rayleigh-Rice distribution is
characterized by the Rice coefficient [21].
Parameters $S_4$, $m$ and the Rice coefficient $c$
subject to the following relations

In accordance with the experimental data
for the normal mid-latitude ionosphere and
for the radio link with the center frequency
$f_0 = 400$ MHz, the flicker index $S_4$ does not
exceed 0.3 ... 0.5, i.e. fading can be attributed
to the class of weak-medium fading, for polar
regions the values of $S_4$ can reach 1 [5] and
in this case fading can be attributed to the
class of strong fading.

Below are the results of estimating
the $S_4$ index based on experimental
measurements of the amplitudes of the signals of the
Cospar-Sarsat satellite system (center frequency $f_0 = 406$ MHz)
and the probability characteristics for
receiving signals calculated using relation
(2) and model descriptions of the
amplitude distribution density $p(A)$ (6),
(7) with parameters corresponding to the
estimates of $S_4$.

4. RESULTS OF CALCULATIONS

The Cospar-Sarsat international satellite
system is designed to determine the location
of emergency beacons operating in the
P-frequency range 406.0 ... 406.1 MHz (uplink)
and in the L-frequency range (downlink) [18].
Information messages of second-generation
beacons are transmitted over a radio link using
digital FM4 signals with an offset [10], a signal
duration of 1 second with a nominal on-time
interval of 30 seconds, and a sync sequence of
160 ms [18]. Relay beacon signals are located...
on global navigation satellites (Glonass, GPS, Galileo [17]).

In Fig. 2 shows fragments of the dependence of the power of the I (t) signals on time, calculated by processing at FIRE them. V.A. Kotelnikov of the RAS of the received signals of the second generation beacons in the form of records in digital format, obtained using the technical means of the receiving station from the GPS mid-orbit navigation satellites (Fig. 2a), Galileo (Fig. 2b) for daytime. Beacon producing countries - France, USA, beacons operate in test mode [18].

The peculiarity of the considered radio link is the joint influence of ionospheric irregularities on the propagation of signals in the P- (uplink) and L- (downlink) frequency ranges.

During processing, point estimates I (t) were calculated for the discrete beacon turn-on time by correlation processing of input realizations with a synchronization sequence.

As a result of processing signals from the Cospas-Sarsat satellite system (time period March-October 2018), a possible range of S4 scintillation index values was determined for the frequency range under study S4 = 0.27 ... 0.49, the approximate fading range of signals amplitude Pf = 5.25 ... 11.20 dB. The corresponding values of the Nakagami distribution parameter (6) are in the range of m = 2.04 ... 3.70, the corresponding range of values of the Rice coefficient of the Rayleigh-Rice distribution (7) c = 4.1 ... 7.4.

In Fig. 3 shows the probabilities of Pb error in the coherent reception of FM4 signals calculated using relations (1), (2) for the scintillation model defined by the Rayleigh-Rice amplitude distribution density (7) with the given estimates of the Rice parameter c.

Curve 1 corresponds to propagation in free space — the probability of Pb = 10–3 is ensured with respect to Eb / N0 = 6.7 dB.

Curve 2 corresponds to the maximum value of the Rice parameter in the range of their estimates c = 7.4, the probability Pb = 10–3 is ensured at the ratio Eb / N0 = 10.0 dB, the energy loss relative to curve 1 reaches 3.3 dB.

Curve 3 corresponds to the minimum value of the Rice parameter in the range of their estimates c = 4.1. It can be seen that the probability of 10–3 is ensured with respect to Eb / N0 = 13.5 dB, which corresponds to the energy loss with respect to curve 1 to 6.8 dB.

Energy losses increase with decreasing Pb error values.

5. CONCLUSION
The descriptions of the models of scintillation (fading) of signals due to random temporal and spatial fluctuations of the electron density of ionospheric inhomogeneities are given. These models fall into two general classes - based on the application of analytical methods for describing signal propagation using the theory of stochastic equations and on the basis of empirical models regarding
the density distribution of signal amplitudes at the input of receivers using a scintillation index (fourth-order statistical moment of signal amplitudes).

The results of the numerical estimation of the scintillation index for the mid-latitude radio line of the Cospas-Sarsat satellite information system are given. A feature of this radio link is the joint influence of ionospheric irregularities on the propagation of signals in the P- (uplink) and L- (downlink) frequency ranges. As a result of signal processing of the Cospas-Sarsat system, a possible range of scintillation index values of 0.27 ... 0.49 was determined. Using the empirical scintillation model, we calculated the probability characteristics of receiving signals with phase shift keying distributed over the radio link under consideration with a given scintillation index range, and determined the required power margin of up to 6.8 dB for the error probability 0.001 with respect to spreading in free space. This must be taken into account when calculating the energy budget of a given radio link.

The study of the statistical characteristics of the studied radio link, in particular, the determination of the time and frequency band of coherence, as well as the specification of the range of the energy reserve, taking into account the noise-resistant coding methods, constitute the direction of prospective studies.

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