Electrochemical Performance of Orthorhombic CsPbI₃ Perovskite in Li-Ion Batteries

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Abstract: A facile solution process was employed to prepare CsPbI₃ as an anode material for Li-ion batteries. Rietveld refinement of the X-ray data confirms the orthorhombic phase of CsPbI₃ at room temperature. As obtained from bond valence calculations, strained bonds between Pb and I are identified within PbI₆ octahedral units. Morphological study shows that the as-prepared δ-CsPbI₃ forms a nanorod-like structure. The XPS analysis confirm the presence of Cs (3d, 4d), Pb (4d, 4f, 5d) and I (3p, 3d, 4d). The lithiation process involves both intercalation and conversion reactions, as confirmed by cyclic voltammetry (CV) and first-principles calculations. Impedance spectroscopy coupled with the distribution function of relaxation times identifies charge transfer processes due to Li metal foil and anode/electrolyte interfaces. An initial discharge capacity of 151 mAh g⁻¹ is found to continuously increase to reach a maximum of ~275 mAh g⁻¹ at 65 cycles, while it drops to ~240 mAh g⁻¹ at 75 cycles and then slowly decreases to 235 mAh g⁻¹ at 100 cycles. Considering the performance and structural integrity during electrochemical performance, δ-CsPbI₃ is a promising material for future Li-ion battery (LIB) application.

1. Introduction

Lithium-ion battery (LIB) technology is considered as the most promising energy storage technology for a wide range of applications in portable electronic devices, electric vehicles (EV) and hybrid electric vehicles (HEV) [1]. In the United States, the Department of Energy has set a cost target for EV battery packs of 125 USD/kWh to be reached in 2022, which was in large part a driving mechanism for LIB technology [2]. Given the projected policy and economic environment, and the intrinsic energy density, long cycle life and stability during long cycling, LIB was made a contending choice as a commercialized energy storage system. Negatively charged lithium anodes play an important role in maintaining the high energy density of the battery [3,4]. In spite of having extremely high theoretical specific capacity (3860 mAh g⁻¹) and low electrochemical potential (~3.04 vs. SHE), the safety risk of metallic Li hinders it from commercialization [5,6]. Commercial LIBs have graphite and lithium titanium oxide (LTO) as anode materials to achieve the market demand [7,8]. However, safety and energy density are a primary concern in those batteries. Graphite anodes face serious stability issues due to the heterogeneous dispersion of conductivity that results in the formation of imperfect solid electrolyte interphase (SEI) [9]. On the other hand, LTO exhibits outstanding stability in long-term cycling at very high C-rate while compromising safety of the battery [10]. However, low energy
density and high operating voltage (1.54 vs. Li/Li⁺) made the commercialization of LTO impractical for a wide range of applications [11,12]. Recently, silicon (Si) has received a lot of attention as a superior anode material due to its high theoretical gravimetric capacity (4200 mA h g⁻¹) and low operation voltage (0.4 V vs. Li/Li⁺) [13]. However, Si experiences 400% volume expansion during charge–discharge process, which affects the long cycle life due to increased probability for mechanical failure [14]. Indeed, given the environmental and material constraints, alternative anode materials for battery research should be developed to identify material choices that exhibit high energy density at low operating voltage while providing outstanding stability and safety. So far, various metal oxides, transition metal sulfides (MoS₂, FeS₂, CoS₈, VS₂), carbon-based materials and Sn-based materials have been investigated as anode materials [15–23].

Organic and inorganic perovskites have a wide range of applications in the fields of solar cell, optoelectronics, laser and electrochromism [24–30]. Previously, scientists worked on the development of organic–inorganic perovskite materials as energy storage materials for Li-ion batteries [31–37]. However, these materials suffer from poor cycling performance and extremely low discharge capacity. The high mobility and Li⁺ storage capability of hybrid perovskites inspire us to use them as active materials. Previously, the orthorhombic CsPbI₃ perovskite was successfully employed on top of Li metal to avoid the deadly dendrite, which causes failures to battery performance [38].

In this work, we report the δ-CsPbI₃ perovskite as an active anode material for Li-ion batteries. We have prepared orthorhombic CsPbI₃ by solution process at a constant temperature of 70 °C. The as-prepared anode material is incorporated in a coin cell against a Li metal. Further, we have investigated the Li storage capability and structural change of host material upon lithiation by the X-ray diffraction and ab initio density functional theory (DFT) calculation. The active material exhibits an initial discharge capacity of 151 mAh g⁻¹ upon 40 mA g⁻¹ current.

2. Experimental Section

2.1. Materials Preparation

The 0.1 M CsPbI₃ solution was prepared by dissolving CsI and PbI₂ (1:1) in a mixture of dimethylformamide (DMF) and dimethyl sulfoxide (DMSO) (vol:1:1) solvents. The as-prepared solution was stirred overnight at 70 °C. Finally, we obtained a yellowish powder which was further ground using mortar and pestle to achieve a fine powder.

2.2. Material Characterization

The morphology of as-prepared active material was investigated using scanning electron microscopy (SEM, JEOL-Japan, JXA-840A). To observe the changes in electrodes after cycling, the coin cells were dissembled and the electrodes were washed using 1,3-dioxolane (DOL). Later, SEM was employed to observe the change in microscopic level. The crystallinity of δ-CsPbI₃ was investigated using X-ray diffraction (XRD) technique using Philips X’Pert diffractometer equipped with a Cu Ka X-ray source (λ = 1.5406 Å) in the 2Ө range of 10°–50°. A JacobsV-670 UV–Vis spectrometer was employed to confirm the presence of absorption peaks of the δ-CsPbI₃ and the band gap was calculated the Tauc plot. XPS (PHI 5000 Versa Probe and an Al Ka X-ray source = 1486.6 eV) was employed to study the binding energy of central metal ion Pb of δ-CsPbI₃. Coin cells were assembled using δ-CsPbI₃ electrode and after first discharge, the cell was opened and the electrode was washed with DOL and dried prior to transfer inside the XPS chamber.

2.3. Electrochemical Measurement

δ-CsPbI₃ loaded electrodes were prepared by mixing 65% δ-CsPbI₃ with 20% carbon black and 15% poly (vinylidene fluoride) (PVDF) binder. The powders were mixed in N-methyl-2-pyrrolidinone (NMP) solvent. As-obtained slurry was coated on Cu foil and dried overnight at 50 °C. Later, the electrode was cut into a 12 mm disk in order to prepare
the coin cells. A 12 mm Li disk was used as the counter electrode. Electrode 1 M LiPF₆ was prepared by mixing ethylene carbonate (EC) and dimethyl carbonate (DMC) (1:1 vol). A pristine Celgard separator was used in between two electrodes for safety reasons. CR2032-type coin cells were assembled in an Ar-filled glove box. Galvanostatic charge-discharge of the coin cells was performed using Think Power battery testing system. PARSTAT MC 1000 electrochemistry workstation was employed to measure the cyclic voltammetry (CV) and electrochemical impedance spectroscopy (EIS). CV tests were performed on the cells in a voltage window of 0.1–3 V and at a scan rate of 0.1 mV s⁻¹. EIS of the cells were performed with an AC amplitude of 10 mV and the frequency range was from 0.01 Hz to 1 MHz.

2.4. Theoretical Treatment

The distribution function of relaxation times (DFRTs) was calculated by the Impedance Spectroscopy Genetic Programming (ISGP) program using the impedance spectra after different cycles [39–44]. As the DFRT approach is applicable only in the Kramers–Krönig (KK) relations compatible regime, the supporting figure (Supplementary Materials, Figure S19) suggests that we can safely identify three peaks within the KK compatible regime (around 0.1 Hz). A similar approach is also adopted elsewhere [39–41,45]. The Rietveld refinement of the XRD pattern was performed by FullProf software package, the Fourier maps were calculated using GFourier and the crystal structure was plotted using VESTA [46,47].

The electronic structure calculation was performed using the Quantum ESPRESSO (QE) distribution, based on density-functional theory, periodic boundary conditions, plane-wave basis sets and pseudopotentials [48,49]. The GBRV ultrasoft pseudopotentials at a plane-wave cutoff of 40 Ry and a charge-density cutoff of 200 Ry were selected for all the calculations with a Gamma-centered k-point mesh [50]. The structure optimization was carried out using a conjugate-gradient algorithm until the forces on each atom were less than 0.01 eV Å⁻¹ and the total energy was converged up to 10⁻⁵ eV. To study the lithium intercalated structures of 0.5, 1, 2, 3 and 4, supercells of LiCsPbI₃, LiCsPbBr₂, LiCsPbI₂, LiCsPbBrI and LiCsPbO₂ structures with their stable configuration have been considered, respectively.

3. Results and Discussion

1 M δ-CsPbI₃ was prepared by mixing CsI and PbI₂ in DMF and DMSO solvent which results in a yellow solution (Figure S1a). The mixture was stirred overnight inside a N₂-filled glove box at 70 °C. We received a yellowish CsPbI₃ powder as our final active material (Figure S1b). Figure 1 exhibits the crystal structure of δ-CsPbI₃ and diffraction peaks assigned to orthorhombic structure or δ-phase as determined by the Rietveld refinement [51]. Although most of the perovskites are air-sensitive, δ-CsPbI₃ is extremely stable in air [52]. δ-CsPbI₃ maintains its crystallinity as confirmed by the present solution method. According to previous reports, CsPbI₃ exists in four crystal phases: α, β, γ and δ [53]. Out of all the studied phases, the δ phase is stable in air, and an orthorhombic structure gives sufficient space for Li intercalation–deintercalation [52]. The diffraction pattern of the δ-CsPbI₃ at room temperature has been indexed with orthorhombic phase (space group Pnna) having lattice parameters \(a = 10.4595(3)\ \text{Å}, \ b = 4.8034(15)\ \text{Å}\) and \(c = 17.7754(4)\ \text{Å}\), \(\alpha = \beta = \gamma = 90°\) with a cell volume of 893.055 Å³ which is in accordance with ICSD File #250744. The convergence of the refinement is achieved with the residuals of \(R_w = 7.87\%\), \(R_p = 5.75\%\) and \(\chi^2 = 7.34\), and the refined structural parameters are shown in Table 1. The XRD patterns at different lithiation potentials, determined by CV measurement shown in Figure S2, have a smaller number of peaks in comparison with the pristine powder. Preferred orientation of the Bragg peaks in 2θ around 43° and 50° (Figure S2), and Figure S3 (red curve), supports the conversion reactions. For instance, the peaks at 42.93°, 50.07°, 27.23° in panel (b), and 42.95°, 50.10° in panel (c) can be attributed to δ-CsPbI₃: 25.88°, 31.06° are associated with PbI₂; and 36.02°, 52.02° in panel (d) are from PbLi, as confirmed from ICSD.
Files #250744, #104762, #68819, #24265, respectively. This suggests the coexistence of both pristine and conversion reacted phases at different potentials. The Rietveld refinements of the XRD patterns for both pristine and after 100 cycles (Figure S4 with orthorhombic phase) suggest an increase in unit cell volume \((896.747 - 893.055 = 3.692 \text{ Å}^3)\) indicating lithiation in the cell as well as shallow charging/discharging processes. Furthermore, the lithiation potential calculation by DFT supports this conjecture (see later section). However, the pristine structure changes due to prolonged cycling as well as reaction times of lithiation (Figures S2 and S3). Additionally, the presence of \(\text{Li}_x\text{PF}_y\) moieties as confirmed by the XPS studies further justifies the lithiation process (see Figures S12 and S13, and Equation (1)).

![Rietveld refinement of the XRD pattern of δ-CsPbI3 at room temperature.](image)

Figure 1. Rietveld refinement of the XRD pattern of δ-CsPbI3 at room temperature.

The perovskite structure of δ-CsPbI3 has distorted PbI6 octahedral with a distortion index of 0.025 Å and two Pb-I1, three Pb-I3 and single Pb-I2 distances of 3.232(7) Å, 3.303(7) Å, 3.409(7) Å and 3.060(6) Å, respectively (Figure S5). The average I-I distances are calculated as 4.188(9) Å to 4.803(15) Å. In this context, the electron density distributions along (110) plane are calculated to find any discrepancy between the observed and calculated structure factors. As observed from Figures S6 and S7, the positive electron densities at two highly localized I1 sites justify the proper fitting of the model. Nevertheless, the contour lines suggest the ionic bonding between I sites. The bond-valence parameters have been calculated using the Zachariasen formula to be 1.923 v.u., 1.029 v.u., 1.081 v.u., 0.928 v.u. and 1.106 v.u. for Pb, I1, I2, I3 and Cs sites, respectively. Interestingly, the variation in the bond-valence parameter of I changes due to different atomic coordinates and occupancies (see Table 1). These deviations from their original valence states indicate the strained bonding inside the polyhedral units, and these have a direct effect on battery performance as observed here and elsewhere [54].
Table 1. At room temperature.

| Label of Atom | Atom | Wyckoff Site | Atomic coordinates | Occupancy |
|---------------|------|--------------|---------------------|-----------|
| Pb            | Pb   | 4c           | 0.16099 0.2500 0.43864 | 1         |
| I1            | I    | 4c           | 0.16396 0.2500 0.00343 | 0.995     |
| I2            | I    | 4c           | 0.29759 0.2500 0.28639 | 0.974     |
| I3            | I    | 4c           | 0.03533 0.2500 0.61563 | 1         |
| Cs            | Cs   | 4c           | 0.41541 0.2500 0.66999 | 1         |

Anisotropic thermal factors for the same elements are shown below (Table 2).

Table 2. Anisotropic thermal factors for the same elements.

| Label of Atom | B11      | B22      | B33      | B12      | B13      | B23      |
|---------------|----------|----------|----------|----------|----------|----------|
| Pb            | 0.00148  | 0.04421  | 0.00015  | 0        | 0.00088  | 0        |
| I1            | 0.00302  | 0.01320  | 0.00154  | 0        | 0.00173  | 0        |
| I2            | -0.00049 | 0.05040  | 0.00014  | 0        | 0.00023  | 0        |
| I3            | 0.00127  | 0.04586  | 0.00086  | -        | 0.00028  | 0        |
| Cs            | -        | -0.07601 | 0.00194  | 0        | 0.00076  | 0        |

Figure 2a,b shows regular and uniform nanorod-like structures with an average length of 5–7 μm and diameter of ~ 400 nm of δ-CsPbI₃. The average particle size is 5.5 μm which is calculated from particle size distribution histogram (Figure S8). An absorption peak is observed between 425–450 nm when UV-vis spectra was collected for δ-CsPbI₃ (Figure S9). Tauc plot shows a band gap of 2.644 eV as obtained from the optical absorption spectrum. DFT-based density of states (DOS) is used to investigate how the lithium ion affects the electronic structure of the δ-CsPbI₃ perovskite. From Figure S10, δ-CsPbI₃ shows a band gap of 2.132 eV similar to that obtained for CsPbBr₃ [55]. Our results suggest that Perdew–Burke–Ernzerhof (PBE) underestimates the experimental band gap up to 21%. The top-view SEM image (Figure 2c) shows that the CsPbI₃ is uniformly distributed. Figure 2d shows ~15 μm electrode coated on the Cu foil. An XPS survey spectrum is performed to confirm the compositions of CsPbI₃, Cs, Pb and I₂ at their corresponding binding energies (Figure 2e).
The impedance measurements were carried out in the pristine cell as well as after several charge–discharge cycles (Figure 3a). We have also plotted the impedance data for a Li‖Li symmetric cell for comparison. It is observed that the overall impedance response contains a single semicircular arc with spike-like extensions at low frequencies. With increasing charge–discharge cycles, the diameter of the semicircular arc decreases and shows an abrupt change after maximum 60 cycling. To understand the electrochemical phenomenon qualitatively, the impedance data have been modeled by ISGP, neglecting the capacitive diffusive regime at low frequencies. The program computes the distribution function of relaxation times which comprises several peaks justifying some electrochemical phenomena. As observed from Figure 3b, the Distribution Function Relaxation Times (DFRTs) for all the conditions show three peaks within the experimental frequency range. A tiny peak at the highest frequencies corresponds to the Ohmic drop due to electrolytic resistances (Peak P1). The peak within the $10^2$ to $10^4$ Hz regime corresponds to lithium metal, which is related to the charge transfer processes along the metal surface (Peak P2) [56,57]. The peak shifts towards low frequencies with initial cycling but vanishes after 100 cycles, indicating loss of the reactivity of the lithium metal. Additionally, the charge transfer resistance decreases with increasing cycles, justifying good electronic transport (Figure S11). At pristine condition, due to the initial SEI layer at anode, lithium activity maintains...
an equilibrium which dilutes with increasing cycle numbers and as a result, the resistance decreases. Secondly, the peak within $10^{-10^{-2}}$ Hz corresponds to the charge transfer of lithium ions at the anode/electrolyte interface and exhibits strong frequency dependence (Peak P3) [56]. This charge transfer process becomes slower with increasing cycle numbers. A close inspection shows that the peak P3 for cycles 2 and 3, unlike cycles 1 and 70 (Figure 3a inset), has broader time distribution which is due to small variation in their impedance plots. It seems that the charge transfer processes are important in explaining the state-of-health of the batteries and can be used as representative fingerprinting for such high-capacity behavior.

![Figure 3](image-url)

**Figure 3.** (a) Nyquist plots of impedance after different cycles along with symmetric Li cell. (b) DFRTs for Li-ion batteries after different cycling measurements. Inset of (a): zoom-in figure of impedance.

The XPS spectra in Figure 4a reveal the electrochemical interaction of lithium with δ-CsPbI$_3$. Both spectra contain peaks from Cs (3d, 4d), Pb (4d, 4f, 5d) and I (3p, 3d, 4d) and with a decrease in their intensities (except O and F) in the material after charging–discharging cycle, suggesting a coating of surface film on the active electrode [58–59]. A tiny peak at 532 eV of O 1s is observed which could be due to Li alkox species: CH$_3$OLi (from reduction of DMC) and (CH$_2$OLi)$_2$ (from reduction of EC) compounds as obtained (Figure 4a) [60]. Secondly, the peak becomes sharper after the discharging condition. A peak at 286 eV corresponds to unavoidable C 1s peak (C-C). XPS spectra exhibit the Pb$^{2+}$ states and the binding energies 138.8 and 143.7 eV, corresponding to 4f$_{7/2}$ and 4f$_{5/2}$, respectively.
Previously, it was reported that due to the influence of external ions/factors, there might be change in oxidation state of Pb$^{2+}$ or redistribution of electron charges around Pb$^{2+}$ [55,61]. The binding energies shifted from 138.8 to 137.7 eV and 143.7 to 142.7 eV, respectively. The absence of I 4d$^{5/2}$ peak at 49 eV and I 3d$^{5/2}$ and 3d$^{3/2}$ around 630 eV confirm that the discharging effect is associated with iodine. Overall, as obtained from the XRD refinement results, the PbI$_6$ octahedra are associated with the lithiation process. As expected, no LiF peaks are detected in the pristine electrode. Another important feature is that the F 1s spectra at 688 eV was observed after discharged (Figure S12). The F 1s spectra consists of two peaks around 685 eV (due to LiF) and 687 eV due to Li$_x$PF$_y$ with the following reduction reactions [60]:

$$LiPF_6 + (5 + x - y)e^- + (5 + x - y)Li^+ \rightarrow (6 - y)LiF + Li_xPF_y$$  (1)

The Li 1s peak at 56 eV can be observed (LiF) (Figure S13) and the P 2p peak (138 eV) is masked with Pb 4f, inferring Li$_x$PF$_y$ moieties. On the other hand, a relatively higher O 1s peak suggests Li$_x$PF$_y$O$_z$ moieties. Since the electrode is washed thoroughly with DMC before the measurements, the peaks relating to F 1s and P 2p may not be due to original LiPF$_6$. Thus, the surface of the electrode is modified by different species such as LiF, Li$_x$PF$_y$ and Li$_x$PF$_y$O$_z$ due to the first discharged cycle.

Figure 4. (a) Full scan XPS spectra of CsPbI$_3$, (b) XPS spectra of Pb 4f before and after electrochemical discharging process.

Figure 5a represents the cyclic voltammetry (CV) of Li-ion coin cells with δ-CsPbI$_3$ as anode materials. The CV is performed within a voltage range of 0.1–3 V at a scan rate of 0.1 mV s$^{-1}$. During delithiation, two peaks at 0.58 V and 0.7 V vs. Li/Li$^+$ are noted, which have been reported for MAPbBr$_3$ [62]. Furthermore, during lithiation, the peaks between 0.4 V–0.6 V can be attributed to Li$_x$Pb such as LiPb and Li$_x$Pb as identified previously for PbO$_2$, PbO and Pb electrodes [63]. It is noted that similar characteristic features of alloying and dealloying have been reported for 2D and 3D hybrid perovskites [64].
defragments due to alloying are also noted from SEM images (Figures S14 and S15). Concerning the peak positions and their variations in CV and XPS spectra, the following electrochemical reaction is proposed: $\text{CsPbI}_3 + \text{Li}^+ + e^- \rightarrow \text{Li}_x \text{CsPbI}_3$ (during discharging). Theoretically, to determine the lithiation potential in CsPbI$_3$, Li-ion is considered to be inserted in the primitive orthorhombic unit cell of CsPbI$_3$ (20 atoms) with its different stoichiometries. Since in the pristine structure, Cs, Pb and I-ions occupy 4c sites, we have chosen 8d, 4c, 4b and 4a as interstitial lithiation sites. The geometry optimization calculations suggest that the ground state configuration energy is the same for all the sites with the same Li atoms per unit cell (as listed in Table S1). Based on the CV curve (Figure 5a), the peaks between 0.4 V–0.6 V can be associated with LiPb such as LiPb, Li$_2$Pb as previously observed for Pb-, PbO$_2$- and PbO-based electrodes, whereas the peak around 1.4 V is due to Li$_x$CsPbI$_3$ (Figure 6a). To identify the quantity of the intercalated lithium ion per unit cell, we have performed geometry optimization in QE with $x = 0.5$, 1, 2, 3 and 4 Li-ions only at 8d site of the host structure. Figure 6a shows a difference in the insertion voltages at each $x$ referring to a single-phase and one step mechanism of lithiation, and a variation in the lattice parameters is observed (Figure 6b). For instance, the lithiation reaction having 0.5 Li-ion will be CsPbI$_3 + 0.5\text{Li}^+ + e^- \rightarrow \text{Li}_{0.5}\text{CsPbI}_3$. Nevertheless, the simulated XRD patterns of Li$_{0.5}$CsPbI$_3$, Li$_2$CsPbI$_3$, Li$_3$CsPbI$_3$ and Li$_4$CsPbI$_3$ are similar to that of the pristine structure (orthorhombic phase). Indeed, the parent structure is an indirect band gap semiconductor, but with an insertion of 0.5 Li per unit cell, the structure becomes semi-metallic (Figure S10). However, the band gap widens with an increase of lithium content (Figure S16). The formation energy and the average Li intercalation potential curves in Figure 6 show that with increasing lithiation, the magnitude of the formation energy increases and the average Li intercalation potential decreases. As the formation energy of lithiation is less than 0 eV, the intercalation is definitely possible. However, the discharge curve in CV suggests that conversion reaction relating LiPb starts around 0.6 V (Figure 5a). Thus, we can confirm that both intercalation and conversion have occurred during the whole discharge process, and due to the conversion process, a structural phase transition may exist (Figures S2 and S3). Comparing Figures 5a,b and 6, the kinks in the discharge CV curve around 1.4 V and 1.0 V can be attributed to Li$_{0.5}$CsPbI$_3$ and Li$_2$CsPbI$_3$, respectively. It is noted that similar insertion and conversion reactions during discharging are reported for the W$_{18}$O$_{39}$ anode [65,66]. Nevertheless, the unit cell volume expands for the lithiated structures (Figure 6b). It is observed that a small volume change is a direct consequence of the low-density perovskite structure of allotrope CsPbI$_3$, facilitating Li-ion movement with minimal lattice distortion. The CV results indicate that the electrochemical reduction of $\delta$-CsPbI$_3$ is reversible. Additionally, the peaks at charging condition do not show any voltage drift, signifying a steady formation of SEI layer as well as no electrolyte oxidation during the cell operation [67]. Figure 5b exhibits $dQ/dV$ profiles ranging from 0.01 to 3 V without any loss of capacitance during the first cycle. This feature establishes its applicability as an anode material in comparison with graphite [68]. The post-cycling cell (after 100 charge–discharge cycles) was tested in similar condition, in order to observe the impact of lithiation/delithiation on the CV of the active material. Figure S17 shows the reversible CV profile of $\delta$-CsPbI$_3$ after 100 cycles inferring negligible loss of active material during long cycling, and as a result, the capacity is maintained around the same as that after 75 cycles. The SEM image of $\delta$-CsPbI$_3$ electrode after 100 cycles observed shows a smooth and integrated surface (Figure S14). Figure S15 is the cross-sectional image of the electrode after 100 charge–discharge cycles showing the coating remains intact with the Cu foil. This study suggests that the $\delta$-CsPbI$_3$ electrode has good integrity throughout a long cycle life and that there is no loss of active materials during cycling. Later, we disassembled the cell and washed and dried the electrode to observe the XRD peak. Figure S4 shows major peaks for $\delta$-CsPbI$_3$ maintained after 100 cycles.
Figure 5. Electrochemical performance of δ-CsPbI₃ active materials in Li-ion half-cell. (a) Cyclic voltammetry trace of a δ-CsPbI₃ half-cell within a voltage window of 0.1–3.0 V, recorded at a scan rate of 0.1 mV s⁻¹. (b) Incremental capacity dQ/dV vs. cell potential as recorded at a rate of 40 mA g⁻¹ for first electrochemical cycle.

Figure 6. (a) Calculation of formation energies as a function of lithium content. The right vertical axis represents the lithiation potential as a function of lithium content in (a). Theoretically calculated volume change (%ν) as a function of lithium content is shown in (b).

Figure 7 exhibits the long-term cycling performance of Li-ion battery using δ-CsPbI₃ as active material at 40 mA g⁻¹. The δ-CsPbI₃ electrode suffers from a capacity drop in the first six charge–discharge cycles from 151 mA h g⁻¹ to 149.8 mA h g⁻¹. From seventh cycle, the capacity started to increase steadily with an increase rate of 2.12 mA h g⁻¹ per cycle to reach a maximum of ~275 mA h g⁻¹ at 65 cycles. After that, it drops to ~240 mA h g⁻¹ at 75 cycles and then slowly decreases to ~235 mA h g⁻¹ at 100 cycles. This phenomenon suggests that Li-ion accessibility increases inside the active material along with oxygen containing functional groups with increasing the cycling number. This can be observed in Figure S9 with reversible CV of active material after 70 cycles. Additionally, this could be due to a large surface area (SEM) and high oxygen adsorption at the surface as supported by the XPS spectrum (Figure 4a). There are a few valid reasons behind a slight increase in discharge capacity which are as follows: (a) due to increase in activation of the active materials, (b) activation of defects in the active materials, and (c) enlargement of active materials’ interatomic space after a few charge–discharge cycles (as understood from increment...
Another possibility is that there is a conversion reaction happening at low voltages which could suppress the intercalation, and due to the conversion reaction, the structural defragment is obtained (Figure S2). Overall, the capacity has increased Li-ion accessibility inside the active material along with oxygen containing functional groups with an increase in the cycling number. The Coulombic efficiency was maintained throughout the cycling performance, which suggests a reversibility of Li insertion/disinsertion during electrochemical process. This cycling performance reveals that Li can intercalate–deintercalate during the charge–discharge process, and it is a reversible process [72,73]. This suggests the Li-ion movement during cycling was smooth and avoided thicker SEI formation [74].

Furthermore, as suggested by the dQ/dV plot (Figure 5b), there is a negligible SEI layer in comparison with that of graphite [74].

**Figure 7.** Long-term cycling performance of a Li-ion battery featuring δ-CsPbI₃ as active materials, tested at 0.1C-rate.

We have shown the cycling performance at different C-rates to evaluate the cycling stability of the active material. As displayed in Figure S18, the cell exhibits an initial discharge capacity of 181 mAh g⁻¹ when discharged at 0.1 C-rate. After tested through harsh charge–discharge conditions at different C-rates, the cell retains ~84% of the initial discharge capacity, which shows outstanding cycling performance of the CsPbI₃ as an active anode material.

### 4. Conclusions

We have successfully introduced δ-CsPbI₃ as an active anode material for Li-ion battery. The perovskite structure of δ-CsPbI₃ has distorted PbI₆ octahedra as calculated by the X-ray refinement. Ionic bonding between I sites are observed along with a partial vacancy up to 0.03%. We propose that the presence of the functional species has a beneficial effect in terms of discharge capacity. Specifically, the charge transfer processes along the Li metal surface and that at the anode/electrolyte interface are identified, and both of them show constant frequency shifts with increasing cycling. We conclude that both intercalation and conversion have occurred during the whole discharge process using ex situ XRD, CV and first-principles calculations. The electrochemical performance of δ-CsPbI₃ shows outstanding potential as a promising anode material for commercialized Li-ion batteries. We need more extensive studies to stabilize performance while at the same time increasing their capacity.

**Supplementary Materials:** The following are available online at www.mdpi.com/10.3390/ma14195718/s1, Figure S1: (a) Photograph of δ-CsPbI₃ solution during preparation, and (b) as-prepared δ-CsPbI₃ powder; Figure S2: XRD of fresh CsPbI₃ and after one cycle at various voltages during discharging; Figure S3: XRD patterns of pristine CsPbI₃, CsPbI₃ with LiPF₆ salt and CsPbI₃ after dipping in electrolyte for different time along with Li metal foil at room temperature; Figure S4: Rietveld refinement of the XRD pattern of δ-CsPbI₃ after 100 cycles at room temperature; Figure S5: Orthorhombic structure of δ-CsPbI₃. PbI₆ octahedral units are shown in gray color; Figure
S6: Electron density distribution of observed structure factors along (110) plane of δ-CsPbI3; Figure S7: Electron density distribution of calculated structure factors along (110) plane of δ-CsPbI3; Figure S8: Particle size distribution for CsPbI3 from SEM image; Figure S9: Optical absorption spectra of CsPbI3 collected by dissolving as-prepared CsPbI3 powder in DMF solvent. (a) UV-Vis spectra and (b) Tauc plot.; Figure S10: Total DOS for CsPbI3 and LiCsPbI3. The vertical line represents Fermi level: EF; Figure S11: Variation of the charge transfer resistances along the metal surface with cycles; Figure S12: XPS spectra of F 1s before and after electrochemical discharging process; Figure S13: XPS spectra of Li 1s after electrochemical discharging process; Figure S14: Top view FESEM image of δ-CsPbI3 electrode after 100 charge-discharge cycles at 40 mA g⁻¹. Figure S15: Cross sectional FESEM image of δ-CsPbI3 electrode after 100 charge-discharge cycles at 40 mA g⁻¹; Figure S16: Total DOS for several lithiated compounds along with pristine CsPbI3. The vertical line represents Fermi level: EF; Figure S17: Cyclic voltammetry trace of a δ-CsPbI3 half-cell after 100 Charge-discharge cycles within a voltage window of 0.1–3.0 V, recorded at a scan rate of 0.01 mV s⁻¹.; Figure S18: Cycling performance of CsPbI3 anode materials at different C-rate.; Figure S19: KK compatibility test.; Table S1: Ground state energy for different configuration of lattice sites with different concentration of lithiation.

Author Contributions: Data curation, N.K., P.-W.C. and Y.-H.S.; formal analysis, N.K., T.P. and P.-W.C.; funding acquisition, M.-K.W.; investigation, M.-K.W. and P.M.W.; methodology, N.K., P.-W.C., Y.-H.S. and A.S.; project administration, M.-K.W. and P.M.W.; resources, M.-K.W.; software, T.P.; supervision, C.-W.C., M.-K.W. and P.M.W.; writing—original draft, N.K.; writing—review & editing, N.K., T.P., M.-K.W. and P.M.W. All authors have read and agreed to the published version of the manuscript.

Funding: The authors greatly appreciate the financial supports from the Executive Yuan through its Forward-Looking Research Grant Number 110-0210-04-20-01. Dr. Phillip M. Wu acknowledges the support of visiting scholar program under Grant Number 109-2811-E-027-506 from the Ministry of Science and Technology (MOST) of Taiwan.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: The data presented in this study are available on request from the corresponding author.

Conflicts of Interest: The authors declare no conflict of interest.

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