Hidden Photon CDM Search at Tokyo

Jun’ya Suzuki, Yoshizumi Inoue, Tomoki Horie, Makoto Minowa
The University of Tokyo, Japan

DOI: http://dx.doi.org/10.3204/DESY-PROC-2015-02/junya.suzuki

We report on a search for hidden photon cold dark matter (HP CDM) using a novel technique with a dish antenna. We constructed two independent apparatuses: one is aiming at the detection of the HP with a mass of $\sim$ eV which employs optical instruments, and the other is for a mass of $\sim 5 \times 10^{-5}$ eV utilizing a commercially available parabolic antenna facing on a plane reflector. From the result of the measurements, we found no evidence for the existence of HP CDM and set upper limits on the photon-HP mixing parameter $\chi$.

1 Introduction

Astronomical observations of the past decades reveal that there exists invisible non-baryonic matter (dark matter, DM) in the universe. Exploring the nature of DM is one of the most important issues in astrophysics and cosmology today, and a variety of experiments have been carried out to directly detect DM particles.

The most prominent candidate for DM is the Weakly Interacting Massive Particle (WIMP), and most of the current experiments aim at the detection of WIMPs. However, there are alternative candidates to account for the features of DM, and Weakly Interacting Slim Particles (WISP), e.g. axion-like particles (ALP) or hidden-sector photons (HP), can be the main component of DM [1].

Hidden photon cold dark matter (CDM) can be experimentally investigated via kinetic mixing ($\chi/2) F_{\mu\nu} \tilde{X}^{\mu\nu}$ between photons and hidden photons. For example, the Axion Dark Matter eXperiment (ADMX) [2], which employs a resonant cavity and magnetic field to search for axion dark matter, also has sensitivity to hidden photon CDM, and its non-detection of the signal [3, 4, 5, 6, 7] was translated to the upper limit for the kinetic mixing parameter $\chi$ [1].

Additionally, a novel method with a spherical mirror to search for HP CDM was recently proposed [8], with which wider mass-range can be probed without rearranging the setup. In this method, ordinary photons of energy $\omega \simeq m_{\gamma'}$ induced by HP CDM via kinetic mixing are emitted in the direction perpendicular to the surface of the mirror, resulting in concentration of the power to the center of the mirror sphere.

This method using a spherical reflector is extremely simple, and can be implemented relatively easily. To confirm its feasibility in real situations, we planned and carried out two experiments to search for HP CDM in two different mass regions: one is for $m_{\gamma'} \sim$ eV using optical equipments and the other for $m_{\gamma'} \sim 50 \mu$eV employing RF instruments. Here we report on the preparations and the results of those searches for HP CDM using the dish method.

Axion–WIMP 2015 145
2 Optical search

For the search in $m_{\gamma'} \sim \text{eV}$, we need a spherical mirror and a photodetector. Non-relativistic HPs near the surface of a reflector induce emission of photons in the direction perpendicular to the surface. A photodetector is placed at the point of convergence and detects emitted photons.

We used a parabolic mirror as a ‘dish’. The parabolic mirror is 500 mm in diameter, 1007 mm focal length and the focal spot diameter is 1.5 mm. We used a parabolic surface instead of a spherical surface originally proposed in Ref. [8] to reuse the mirror which had been employed in the solar HP helioscope [9]. From the diameter and the focal length of the parabolic mirror, photons emitted perpendicularly to the surface are calculated to concentrate to a small area of 4 mm in diameter at twice the focal length of the mirror, which is small enough compared to the effective area of the photodetector.

A photomultiplier tube (PMT) was employed as the detector of emitted photons. We selected Hamamatsu Photonics R3550P because of its low dark count rate of $\sim 5\text{Hz}$. We used a motorized stage to shift the position of the PMT, which enabled us to measure background noise.

The mirror and the detector were mounted on a steel frame, which rigidly holds the arrangement (Fig. 1). After installing the optical equipments, this frame was wrapped with black polyethylene sheets to shield from ambient light. Additionally, the whole setup was installed in a light-tight box of $1\times1\times3\text{m}$ to attain higher light-tightness.

With this setup, we carried out the experimental search for HP CDM in the eV mass range [10]. The overall duration of the measurement was $8.3 \times 10^5\text{s}$ for each configuration: with the PMT at the position of convergence of the HP CDM signal (signal, S) and at the position displaced by 25 mm from position S (background, B). We found no excess in count rate measured at position S compared to at position B. We translated this non-detection result to the limit for the mixing parameter $\chi$ (Fig. 2).

3 RF search

We also targeted detection in $K_u$ band ($\sim 12\text{GHz}$) for the feasibility test of the ‘dish’ method. We can use commercially available dish antennas for this frequency region, though they usually have parabolic shape, which cannot be approximated as spherical shape because of their short focal lengths compared to their diameters. In order to overcome this problem, we let our dish face a plane reflector, from which plane radio waves of HP CDM origin would be emitted perpendicularly to the surface. Because parabolic dishes concentrate plane waves to their focal point, the amplification of HP CDM signal properly works. We used an Anstel lar SXT-220 as a dish, which is 2.2 m in diameter and designed for CS broadcast reception. A huge plane reflector was constructed by combining four aluminum plates on a rigid frame.
Figure 2: Preliminary results of the experimental searches for HP CDM. The vertical axis shows the mixing parameter $\chi$, and the horizontal axis shows the mass of hidden-photon $m_{\gamma'}$. The red colored regions are excluded by our results for two experimental setups. With the optical setup, we excluded the area around $m_{\gamma'} \sim eV$. The search in K_u band excluded the region around $m_{\gamma'} \sim 50 \mu eV$. For a descriptions of the other colored areas, see Ref. [10].

For the converter, we selected Norsat 4506B, which down-converts the signal with the local frequency of 11 GHz. The output of the converter was connected to the Fast Fourier Transform (FFT) analyzer, Rohde & Schwarz FSV-4. The signal of the existence of HP CDM would be seen as a spectral line with a broadening of $\Delta f/f \sim 10^{-6}$ due to the velocity dispersion of DM.

After the calibration, the setup for the experimental search was constructed by setting the dish in front of the plane reflector (Fig. 3).

Using this setup, we actually carried out the experimental search for four days. We observed no signal-like excess in the power spectrum and set an upper limit for the parameter $\chi$ (Fig. 2). Although the limit is narrow in the sensitive mass region, we can expand it only by replacing the converter for one which is capable of handling wider frequency range.
4 Conclusion

We constructed two apparatuses utilizing a novel method using a dish antenna. One uses an optical mirror for the survey in $m_{\gamma'} \sim eV$, and the other uses a dish antenna for CS broadcast reception to search HPs with $m_{\gamma'} \sim 50 \mu eV$. We actually carried out the experimental search, and found no evidence for the existence of HP CDM. From the result, we set upper limits on the photon-HP mixing parameter $\chi$ in two different mass regions (Fig. 2).

Acknowledgments

T. Horie acknowledges support by Advanced Leading Graduate Course for Photon Science (ALPS) at the University of Tokyo. This research is supported by the Grant-in-Aid for Challenging Exploratory Research by MEXT, Japan, and also by the Research Center for the Early Universe, School of Science, the University of Tokyo.

References

[1] P. Arias, D. Cadamuro, M. Goodsell, J. Jaeckel, J. Redondo et al., JCAP 06 013 (2012).
[2] H. Peng, S. Asztalos, E. Daw, N. A. Golubev, C. Hagmann et al., Nucl. Instr. Meth. A 444 569 (1999).
[3] S. DePanfilis, A. C. Melissinos, B. E. Moskowitz, J. T. Rogers, Y. K. Semertzidis et al., Phys. Rev. Lett. 59 839 (1987).
[4] W. U. Wuensch, S. De Panfilis-Wuensch, Y. K. Semertzidis, J. T. Rogers, A. C. Melissinos et al., Phys. Rev. D 40 3153 (1989).
[5] C. Hagmann, P. Sikivie, N. S. Sullivan, and D. B. Tanner, Phys. Rev. D 42 1297(R) (1990).
[6] S. Asztalos, E. Daw, H. Peng, L. J Rosenberg, C. Hagmann et al., Phys. Rev. D 64 092003 (2001).
[7] S. J. Asztalos, G. Carosi, C. Hagmann, D. Kinion, K. van Bibber et al., Phys. Rev. Lett. 104 041301 (2010).
[8] D. Horns, J. Jaeckel, A. Lindner, A. Lobanov, J. Redondo, and A. Ringwald, JCAP 04 016 (2013).
[9] T. Mizumoto et al., JCAP 07 013 (2013) [arXiv:1302.1000 [hep-ex]].
[10] J. Suzuki, T. Horie, Y. Inoue, and M. Minowa, arXiv:1504.00118 [hep-ex]. Accepted for publication in JCAP.