GENERALIZED WRIGHT FUNCTION AND ITS PROPERTIES USING EXTENDED BETA FUNCTION

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Abstract. In this paper, we introduce a new generalization of the Wright function by using an extended beta function and study some classical properties of this function. We establish several formulas involving integral transforms (e.g. Jacobi transform, Gegenbauer transform) and the generalized family of Wright function that does not seem to be reported in the literature even for the basic Wright function. Furthermore, we discuss other results including the recurrence relation, derivative formula, fractional derivative formula and also a partly bilateral and partly unilateral generating relation for the generalized Wright function.

1. Introduction and preliminaries

The special functions of mathematical physics are found to be very useful in finding the solutions of initial and boundary value problems associated with partial differential equations and fractional differential equations. Special functions have widespread applications in areas of engineering as well, and often new perspectives in special functions are motivated by such connections. Several special functions, called now, special functions of fractional calculus play a very interesting role in the solution of fractional order differential equations such as Mittag-Leffler function, Wright function, and its auxiliary functions and Fox $H$- function. The Wright function was first introduced $[21, 22, 23, 24]$ during the study of the asymptotic theory of partitions of natural numbers. It plays an important role in the solution of linear partial differential equations and subsequently, many other applications for instance, in the Mikusiński's operational calculus ([7], see also [8]) and in the theory of Hankel type integral transforms have been found. Recently, this function has appeared in the area related to the partial differential equation of fractional order. Considering the boundary value problems for the fractional diffusion-wave equation, that is, the linear partial integro-differential equation

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obtained from the classical diffusion or wave equation by replacing the first or second-order time derivatives by a fractional derivative of order \( \alpha \) with \( 0 < \alpha < 2 \), it was found that the corresponding functions can be represented in terms of the Wright function. In this sequel, we intend to introduce a new generalization of Wright function by using the known extended beta function.

Let \( \mathbb{C}, \mathbb{R}^+, \mathbb{N}, \) and \( \mathbb{Z}_0^- \) be the sets of complex numbers, positive real numbers, positive integers, and non-positive integers, respectively, and let \( \mathbb{R}_0^+ := \mathbb{R}^+ \cup \{0\} \) and \( \mathbb{N}_0 := \mathbb{N} \cup \{0\} \).

We recall the classical Wright function \( W_{\alpha,\beta}(z) \) [8, 14] defined by the series representation

\[
W_{\alpha,\beta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{n! \Gamma(\alpha n + \beta)} \quad (\beta \in \mathbb{C}, \alpha > -1).
\]  

(1.1)

The two supplementary functions ([5]; see also [8, 14, 15]) of the above function are given as

\[
M_{\alpha}(z) = W_{-\alpha,1-\alpha}(-z) = \sum_{n=0}^{\infty} \frac{(-1)^n z^n}{n! \Gamma(1 - \alpha(n+1))} \quad (0 < \alpha < 1),
\]

(1.2)

\[
F_{\alpha}(z) = W_{-\alpha,0}(-z) = \sum_{n=1}^{\infty} \frac{(-1)^n z^n}{\Gamma(-\alpha n) n!} \quad (0 < \alpha < 1).
\]

(1.3)

Recently, for real \( \alpha \) and \( \beta, \gamma, \delta \in \mathbb{C}; \alpha > -1, \delta \neq 0, -1, -2, \ldots \) with \( z \in \mathbb{C} \) and \( |z| < 1 \) with \( \alpha = -1 \), Moustafa et al. [5] defined the generalization of the Wright function as

\[
W_{\gamma,\delta}^{\alpha,\beta}(z) = \sum_{n=0}^{\infty} \frac{z^n (\gamma)_n}{n! (\delta)_n \Gamma(\alpha n + \beta)},
\]

(1.4)

where

\[
\frac{\Gamma(\gamma+n)}{\Gamma(\gamma)} = (\gamma)_n = \gamma(\gamma+1)(\gamma+2) \cdots (\gamma+n-1).
\]

is a pochhammer symbol (see, [18, p.2 and pp.4-6]) and \( \Gamma(\cdot) \) is the gamma function (see, [18, Section 1.1]). The two auxiliary functions (generalized Wright type) for all complex variable \( z \neq 0 \) and for any order \( \alpha \in (0, 1) \) were defined as

\[
M_{\alpha}^{\gamma,\delta}(z) = W_{-\alpha,1-\alpha}^{\gamma,\delta}(-z) = \sum_{n=0}^{\infty} \frac{(-1)^n z^n (\gamma)_n}{n! \Gamma(1 - \alpha(n+1)) (\delta)_n},
\]

(1.5)

\[
F_{\alpha}^{\gamma,\delta}(z) = W_{-\alpha,0}^{\gamma,\delta}(-z) = \sum_{n=1}^{\infty} \frac{(\gamma)_n (-1)^n z^n}{(\delta)_n \Gamma(-\alpha n) n!}.
\]

(1.6)

An extension of Euler’s beta function ([2]; see also, [13, p.113(1.2)]) introduced by Chaudhry et al. is defined as

\[
B_p(x, y) = \int_0^1 (1 - t)^{y-1} t^{x-1} e^{-\frac{t^p}{\mu}} \, dt
\]

(1.7)

\[\text{min}(\Re(p), \Re(x), \Re(y)) > 0,\]

and \( B(x, y) (p = 0) \) is the well-known beta function (see, [18, Section 1.1]).
Further, a more generalized form of extended beta function [16] is defined as

\[
B_p^{c,d}(x, y) = \int_0^1 (1-t)^{y-1} t^{x-1} \left( 1 - \frac{p}{1-t} \right) \left[ c + \frac{p}{1-t} \right] \frac{d}{dt} \left[ 1 - \frac{p}{1-t} \right] dt
\]

(1.8)

\((\min(\Re(y), \Re(x)) > 0; \ p \in \mathbb{R}_0^+).\)

Srivastava et al. [19] presented the following generalized beta function

\[
B^{(\nu_m)_{m \in \mathbb{N}_0}}(x, y; p) = \int_0^1 (1-t)^{y-1} t^{x-1} \left( 1 - \frac{p}{1-t} \right) \left[ \nu_m \right]_{m \in \mathbb{N}_0} \frac{d}{dt} \left( 1 - \frac{p}{1-t} \right) dt
\]

(1.9)

\((\min(\Re(y), \Re(x)) > 0; \ p \in \mathbb{R}_0^+),\)

where, \(\mathcal{F}(\nu_m)_{m \in \mathbb{N}_0}; z)\) is the function, analytic within the disk \(|z| < \mathcal{R}\) \((0 < \mathcal{R} < \infty)\) and its Taylor–Maclaurin coefficients be explicitly denoted by the sequence \(\nu_m\).

The classical Jacobi polynomials \(P_n^{(\alpha, \beta)}(z)\) (see, e.g., [4, Chapter 10] and [20, Chapter 4]) in terms of hypergeometric function are defined as

\[
P_n^{(\alpha, \beta)}(z) = \left( \begin{array}{c} \alpha + n \\ n \end{array} \right)_2 F_1 \left[ -n, \frac{\alpha + \beta + n + 1}{2}, 1 - \frac{1-z}{2}; \alpha + 1 \right]
\]

(1.10)

where, \(2 F_1\) is the familiar Gauss hypergeometric function.

The Jacobi transform formula for a function \(f(z)\) is defined as

\[
\mathcal{J}^{(\alpha, \beta)}[f(z); n] = \int_{-1}^1 (1-z)^{\alpha-1} (1+z)^{\beta-1} P_n^{(\alpha, \beta)}(z) f(z) \, dz
\]

(1.11)

\((\min(\Re(\alpha), \Re(\beta)) > -1; \ n \in \mathbb{N}_0),\)

on condition that the integral in (1.11) exists.

Now, on using the definition (1.10), (1.11) can be written as

\[
\int_{-1}^1 (1-z)^{\xi-1} (1+z)^{\eta-1} P_n^{(\alpha, \beta)}(z) \, dz = 2^{\xi + \eta - 1} \left( \begin{array}{c} \alpha + n \\ n \end{array} \right)_3 F_2 \left[ -n, \frac{\alpha + \beta + n + 1, \xi}{\alpha + 1, \xi + \eta}; 1 \right]
\]

(1.12)

\((\min(\Re(\xi), \Re(\eta)) > 0),\)

Lately, Srivastava et al. [17] introduced an interesting Jacobi transform formula

\[
\mathcal{J}^{(\alpha, \beta)}[z^{\rho-1}; n] = \int_{-1}^1 (1-z)^{\alpha} (1+z)^{\beta} P_n^{(\alpha, \beta)}(z) z^{\rho-1} \, dz
\]

\[= 2^{\alpha + \beta + 1} \left( \begin{array}{c} \alpha + n \\ n \end{array} \right)_B(\alpha + 1, \beta + 1)\]
and \( \alpha \beta + 2: \alpha + 1 \), \(-\ldots-\); \(1, 2\) respectively. In view of above relationships, the Gegenbauer transform and Legendre transform other polynomials being the Legendre polynomials \(P_n(z)\), the Chebyshev polynomials \(T_n(z)\) and \(U_n(z)\) (first and second kind), the Gegenbauer polynomials \(C_n^\nu(z)\). We have indeed

\[ C_n^\nu(z) = \left( \frac{v + n - \frac{1}{2}}{n} \right)^{\frac{1}{2}} \left( \frac{2v + n - 1}{n} \right) P_n^{\nu - \frac{1}{2}, \nu - \frac{1}{2}}(z) \]  

(1.14)

and

\[ P_n(z) = C_n^{\frac{1}{2}}(z) = P_n^{(0,0)}(z), \]  

(1.15)

respectively. In view of above relationships, the Gegenbauer transform \(G^\nu[f(z); n]\) and the Legendre transform \(L[f(z); n]\) (see [17]) can be defined as:

\[ G^\nu[f(z); n] = \left( \frac{v + n - \frac{1}{2}}{n} \right)^{\frac{1}{2}} \left( \frac{2v + n - 1}{n} \right) \int_0^1 (1 - z^2)^{\nu - \frac{1}{2}} C_n^{(\nu)}(z) f(z) dz \]  

\[ = \int_{-1}^1 (1 - z^2)^{\nu - \frac{1}{2}} C_n^{(\nu)}(z) f(z) dz \quad \left( \Re(\nu) > -\frac{1}{2}; n \in \mathbb{N}_0 \right) \]  

(1.16)

and

\[ L[f(z); n] = G^{\frac{1}{2}}[f(z); n] = \int_{-1}^1 P_n(z) f(z) dz \quad (n \in \mathbb{N}_0). \]  

(1.17)

**Definition 1.1.** Using the definition in (1.9), we now define a new generalization of the Wright function as follows:

\[ W_{a, b, \gamma, \delta, \lambda, \alpha}^{(\nu)}(z; p) = \sum_{n=0}^{\infty} \frac{B_p^{(\nu),\text{mech}}(\gamma, \lambda) (\alpha n + \beta)}{n! \Gamma(\alpha n + \beta)} \frac{z^n}{(\delta)_n}, \]  

\[ (\alpha = -1, \delta \neq -1, -2, \ldots; \gamma, \delta, \lambda, \alpha, \beta, \sigma, \alpha \in \mathbb{C} \text{ with } z \in \mathbb{C} \text{ and } |z| < 1). \]  

(1.18)

We also introduce the two Wright type auxiliary functions, respectively as

\[ M_{a, b, \gamma, \delta, \lambda, \alpha}^{(\nu),\text{mech}}(z; p) = W_{a, 1-a, \gamma, \delta, \lambda, \alpha}^{(\nu),\text{mech}}(-z; p) = \sum_{n=0}^{\infty} \frac{(-1)^n z^n}{(\delta)_n} \frac{(\gamma)_n}{n! \Gamma(1-a n + \alpha) (\lambda)_n}, \]  

(1.19)

and

\[ F_{a, b, \gamma, \delta, \lambda, \alpha}^{(\nu),\text{mech}}(z; p) = W_{a, b, \gamma, \delta, \lambda, \alpha}^{(\nu),\text{mech}}(-z; p) = \sum_{n=0}^{\infty} \frac{(-1)^n z^n}{(\delta)_n} \frac{(\gamma)_n}{n! \Gamma(-a n) (\lambda)_n}. \]  

(1.20)
Proof. Using (1.9) in (1.18), we have
\[
W_{\{v_{m}\}_{m\in\mathbb{N}_0}}(z; p) = \sum_{n=0}^{\infty} \left\{ \int_{0}^{1} (1 - t)^{\lambda - 1} t^{r + n - 1} \mathcal{F} \left[ \{v_{m}\}_{m\in\mathbb{N}_0}; - \frac{p}{(1 - t)} \right] \right\} \frac{(\sigma)_n}{(\delta)_n B(\gamma, \lambda - \gamma) \Gamma(\alpha n + \beta) n!} z^n.
\]

On interchanging the order of integration and summation in the equation above, which is proven under the given conditions here, we get
\[
W_{\{v_{m}\}_{m\in\mathbb{N}_0}}(z; p) = \int_{0}^{1} t^{r - 1} (1 - t)^{\lambda - 1} \mathcal{F} \left[ \{v_{m}\}_{m\in\mathbb{N}_0}; - \frac{p}{(1 - t)} \right] \sum_{n=0}^{\infty} \frac{(t z)^n}{n! \Gamma(\alpha n + \beta)} \frac{(\sigma)_n}{(\delta)_n B(\gamma, \lambda - \gamma)} d t.
\]

Using (1.4) in the precedent equation, we get the desired result, and in view of (1.19) and (1.20), one can consequently get
\[
M_{\{v_{m}\}_{m\in\mathbb{N}_0}}(z; p) = \frac{1}{B(\gamma, \lambda - \gamma)} \times \int_{0}^{1} (1 - t)^{\lambda - 1} t^{r - 1} \mathcal{F} \left[ \{v_{m}\}_{m\in\mathbb{N}_0}; - \frac{p}{(1 - t)} \right] M_{\alpha, \beta}^{\sigma, \delta}(-t z) d t \quad (2.2)
\]
and
\[
F_{\{v_{m}\}_{m\in\mathbb{N}_0}}(z; p) = \frac{1}{B(\gamma, \lambda - \gamma)} \times \int_{0}^{1} (1 - t)^{\lambda - 1} t^{r - 1} \mathcal{F} \left[ \{v_{m}\}_{m\in\mathbb{N}_0}; - \frac{p}{(1 - t)} \right] F_{\alpha, \beta}^{\sigma, \delta}(-t z) d t. \quad (2.3)
\]

Now, letting \( t = \frac{u}{1+u} \) and \( t = \sin^2 \theta \) in (2.1), we come up with two interesting integral formulas for the extended Wright function in (1.18), which are given in the respective Corollaries.
Corollary 2.2. Let \( c, \gamma, \beta, \alpha, \delta, \lambda, \sigma \in \mathbb{C} \) with \( \Re(\delta) > \Re(\gamma) > 0 \) and \( \Re(\alpha) > 0 \). Also let \( p \in \mathbb{C} \) and \( d \in \mathbb{C} \setminus \mathbb{Z}_0^+ \). Then

\[
W_{a,\beta,\gamma,\delta,\lambda,\sigma}^{(\{v_m\}_{m\in\mathbb{N}_0})}(z; p) = \frac{1}{B(\gamma, \lambda - \gamma)} \int_0^\infty \frac{u^{\gamma-1}}{(1+u)\lambda} \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -\frac{p(1+u)^2}{u}\right] W_{a,\beta}^{\sigma,\delta}(\frac{uz}{1+u}) \, du. \tag{2.4}
\]

Corollary 2.3. Let \( \gamma, \alpha, \beta, \delta, \lambda, \sigma \in \mathbb{C} \) with \( \Re(\delta) > \Re(\gamma) > 0 \) and \( \Re(\alpha) > 0 \). Also, let \( p \in \mathbb{R}_0^+ \). Then

\[
W_{a,\beta,\gamma,\delta,\lambda,\sigma}^{(\{v_m\}_{m\in\mathbb{N}_0})}(z; p) = \frac{2}{B(\gamma, \lambda - \gamma)} \times \int_0^\frac{\pi}{2} \sin^{\gamma-1}\theta \cos^{2\lambda-2\gamma-1}\theta \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -\frac{p}{\sin^2\theta\cos^2\theta}\right] W_{a,\beta}^{\sigma,\delta}(z\sin^2\theta) \, d\theta. \tag{2.5}
\]

Theorem 2.4. The Mellin transform of the generalized Wright function is given by

\[
\mathfrak{M}\left\{ W_{a,\beta,\gamma,\delta,\lambda,\sigma}^{(\{v_m\}_{m\in\mathbb{N}_0})}(z; p); s \right\} = \int_0^\infty p^{s-1} W_{a,\beta,\gamma,\delta,\lambda,\sigma}^{(\{v_m\}_{m\in\mathbb{N}_0})}(z; p) \, dp
\]

\[
= \frac{\Gamma_p(s)\Gamma(\lambda + s - \gamma)\Gamma(\delta)\Gamma(\lambda)}{\Gamma(\gamma)\Gamma(\sigma)\Gamma(\lambda - \gamma)} 2^{\Psi_3}(\sigma, 1, (\gamma + s, 1); (\delta, 1, (\beta, \alpha), (\lambda + 2s, 1); z). \tag{2.6}
\]

Proof. Applying (2.1) to (2.6), we have

\[
\mathfrak{M}\left\{ W_{a,\beta,\gamma,\delta,\lambda,\sigma}^{(\{v_m\}_{m\in\mathbb{N}_0})}(z; p); s \right\} = \frac{1}{B(\gamma, \lambda - \gamma)} \times \int_0^\infty p^{s-1} \times \int_0^1 t^{\gamma-1}(1-t)^{\lambda-\gamma-1} \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -\frac{p}{(1-t)t}\right] W_{a,\beta}^{\sigma,\delta}(t z) \, dt \, dp. \tag{2.7}
\]

On interchanging the order of integration in (2.7), guaranteed under the conditions here, we get

\[
\mathfrak{M}\left\{ W_{a,\beta,\gamma,\delta,\lambda,\sigma}^{(\{v_m\}_{m\in\mathbb{N}_0})}(z; p); s \right\} = \frac{1}{B(\gamma, \lambda - \gamma)} \times \int_0^1 \left\{ t^{\gamma-1}(1-t)^{\lambda-\gamma-1} W_{a,\beta}^{\sigma,\delta}(t z) \right\} \times \int_0^\infty p^{s-1} \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -\frac{p}{(1-t)t}\right] \, dp \, dt. \tag{2.8}
\]

Substituting \( u = \frac{p}{(1-t)t} \) in the inner integral in (2.8), we get

\[
\int_0^\infty p^{s-1} \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -u\right] \, dp = \int_0^\infty u^{s-1} t^{s-1} \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -\frac{p}{(1-t)t}\right] \, du \tag{2.9}
\]

\[
= t^{s}(1-t)^s \int_0^\infty u^{s-1} \mathcal{F}\left[\{v_m\}_{m\in\mathbb{N}_0}; -u\right] \, du
\]

\[
= t^{s}(1-t)^s \Gamma_p^{(\{v_m\}_{m\in\mathbb{N}_0})}(s), \tag{2.10}
\]

\[
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where \( \Gamma_p^{(v_m \in \mathbb{N}_0)}(s) \) is the extended gamma function [17] defined as
\[
\Gamma_p^{(v_m \in \mathbb{N}_0)}(s) = \int_0^\infty t^{s-1} \mathcal{F} \left[ [v_m]_{m \in \mathbb{N}_0}; -t - \frac{v}{t} \right] dt.
\]

Using (2.10) and (2.8), we find
\[
\mathfrak{M} \left\{ W_{a, \beta, \gamma, \delta, \lambda, \sigma}^{(v_m \in \mathbb{N}_0)}(z; p); s \right\} = \frac{\Gamma_p^{(v_m \in \mathbb{N}_0)}(s)}{B(\gamma, \lambda - \gamma)} \int_0^1 t^{r+n+s-1}(1-t)^{\lambda+s-\gamma-1} W_{a, \beta}^{\sigma, \delta}(t \gamma) dt
\]
\[
= \frac{\Gamma_p^{(v_m \in \mathbb{N}_0)}(s)}{B(\gamma, \lambda - \gamma)} \sum_{n=0}^\infty \frac{(\sigma)_n}{(\delta)_n} \frac{z^n}{n!(\alpha n + \beta)} \int_0^1 t^{r+n+s-1}(1-t)^{\lambda+s-\gamma-1} dt
\]
\[
= \frac{\Gamma_p^{(v_m \in \mathbb{N}_0)}(s)}{B(\gamma, \lambda - \gamma)} \sum_{n=0}^\infty \frac{(\sigma)_n}{(\delta)_n} \frac{z^n}{n!(\alpha n + \beta)} \Gamma(\lambda + s - \gamma) \Gamma(\gamma + n + s) \Gamma(\lambda + n + 2s) \Gamma(\lambda - \gamma) \Gamma(\gamma) \Gamma(\sigma) \Gamma(\delta) \Gamma(n)
\]
\[
= \frac{\Gamma(\lambda + s - \gamma) \Gamma(\gamma) \Gamma(\lambda) \Gamma(\sigma) \Gamma(\delta) \Gamma(n) \Gamma(\alpha n + \beta) \Gamma(\delta + n) \Gamma(\sigma) \Gamma(\delta) \Gamma(n)!}{\Gamma(\lambda - \gamma) \Gamma(\gamma) \Gamma(\lambda) \Gamma(\sigma) \Gamma(\delta) \Gamma(n) \Gamma(\alpha n + \beta) \Gamma(\delta + n) \Gamma(\lambda + n + 2s) \Gamma(\lambda + n + 2s) \Gamma(\lambda - \gamma) \Gamma(\gamma) \Gamma(\sigma) \Gamma(\delta) \Gamma(n)!}
\]
which in terms of Fox-Wright function \( p \psi_q \), yields
\[
\mathfrak{M} \left\{ W_{a, \beta, \gamma, \delta, \lambda, \sigma}^{(v_m \in \mathbb{N}_0)}(z; p); s \right\} = \frac{\Gamma_p^{(v_m \in \mathbb{N}_0)}(s)}{\Gamma(\gamma) \Gamma(\lambda) \Gamma(\delta) \Gamma(\sigma) \Gamma(\lambda - \gamma) \Gamma(\gamma) \Gamma(\sigma) \Gamma(\lambda - \gamma)} 2 \psi_3 \left[ \begin{array}{c} \sigma, 1, (\gamma + s, 1); \\ \delta, 1, (\beta, \alpha), (\lambda + 2s, 1); \\ z \end{array} \right].
\]

**Corollary 2.5.** Taking \( s = 1 \) in Theorem 2.4 and using \( \Gamma^{(c, d)} \) (1) = \( \frac{\Gamma(d) \Gamma(c - 1)}{\Gamma(c)} \), we get
\[
\mathfrak{M} \left\{ W_{a, \beta, \gamma, \delta, \lambda, \sigma}^{(v_m \in \mathbb{N}_0)}(z; p); 1 \right\} = \frac{\Gamma(\delta) \Gamma(\lambda) \Gamma(\delta + 1) \Gamma(\sigma) \Gamma(\lambda + 1 - \gamma) \Gamma(\alpha n + \beta) \Gamma(\delta + n) \Gamma(\lambda + n + 2s) \Gamma(\lambda + n + 2s)}{\Gamma(\gamma) \Gamma(\sigma) \Gamma(\delta) \Gamma(\sigma) \Gamma(\gamma) \Gamma(\lambda - \gamma)} 2 \psi_3 \left[ \begin{array}{c} \sigma, 1, (\gamma + 1, 1); \\ \delta, 1, (\beta, \alpha), (\lambda + 2, 1); \\ z \end{array} \right].
\]

**Corollary 2.6.** By the Mellin inversion formula, we have the following complex integral representation for \( W_{a, \beta, \gamma, \delta, \lambda, \sigma}^{(v_m \in \mathbb{N}_0)}(z; p); z \):
\[
W_{a, \beta, \gamma, \delta, \lambda, \sigma}^{(v_m \in \mathbb{N}_0)}(z; p) = \frac{1}{2\pi i} \int_{-\infty}^{i\infty} \Gamma_p^{(v_m \in \mathbb{N}_0)}(s) \Gamma(\lambda + s - \gamma) \Gamma(\delta) \Gamma(\lambda) \Gamma(\gamma) \Gamma(\sigma) \Gamma(\lambda - \gamma)
\]
\[
\times z^\psi_3 \left[ \begin{array}{c} \sigma, 1, (\gamma + s, 1); \\ \delta, 1, (\beta, \alpha), (\lambda + 2s, 1); \\ z \end{array} \right] p^{s-1} ds.
\]
We present a recurrence relation for (1.18) given by the undermentioned corollary:

**Corollary 2.7.** Let \( \gamma, \delta, \lambda, \sigma, \alpha, \beta \in \mathbb{C} \) with \( \Re(\delta) > \Re(\gamma) > 0 \) and \( \Re(\alpha) > 0 \). Also, let \( p \in \mathbb{R}^*_0 \). Then

\[
W_{\alpha, \beta, \gamma, \delta, \lambda, \sigma}^{(\nu_m \text{me}_N)}(z; p) = \beta W_{\alpha, \beta+1, \gamma, \delta, \lambda, \sigma}^{(\nu_m \text{me}_N)}(z; p) + \frac{\alpha \sigma \gamma z}{\lambda \delta} W_{\alpha, \alpha+1, \gamma+1, \delta+1, \lambda+1, \sigma+1}^{(\nu_m \text{me}_N)}(z; p). \tag{2.14}
\]

**Proof.** By using (3.12), it is easy to obtain the desired result. The detailed proof is omitted.

**Theorem 2.8.** With the conditions already mentioned for (1.18), the following Jacobi transform is true:

\[
J^{\alpha, \beta} \left[ z^{p-1} W_{\alpha, \beta, \gamma, \delta, \lambda, \sigma}^{(\nu_m \text{me}_N)}(xz; p); n \right] = 2^{\alpha + \beta + 1} \left( \frac{\alpha + n}{n} \right) B(\alpha + 1, \beta + 1) \sum_{k=0}^{\infty} \mathcal{U}_k \tag{2.15}
\]

where, the coefficients \( \mathcal{U}_k \) are given by

\[
\mathcal{U}_k = \sum_{k=0}^{\infty} \frac{B_p^{(\nu_m \text{me}_N)}(\gamma + k, \lambda - \gamma)}{B(\gamma, \lambda - \gamma) \Gamma(\alpha k + \beta) \Gamma(\delta k)} \tag{2.16}
\]

and we assume that the Jacobi transform in (2.15) exists.

**Proof.** Applying the definition (1.11) in concurrence with (1.18), we have

\[
J^{\alpha, \beta} \left[ z^{p-1} W_{\alpha, \beta, \gamma, \delta, \lambda, \sigma}^{(\nu_m \text{me}_N)}(xz; p); n \right] = \int_{z=1}^{z=1} z^{p-1} (1+z)^\beta (1-z)^\alpha P_n^{\alpha, \beta}(z) W_{\alpha, \beta, \gamma, \delta, \lambda, \sigma}^{(\nu_m \text{me}_N)}(xz; p) dz
\]

\[
= \int_{z=1}^{z=1} z^{p-1} (1+z)^\beta (1-z)^\alpha P_n^{\alpha, \beta}(z) \left( \sum_{k=0}^{\infty} \mathcal{U}_k \frac{(xz)^k}{k!} \right) dz, \tag{2.17}
\]

where \( \mathcal{U}_k \) are given by (2.16). Now, by changing the summation and integration order, and applying the Jacobi transform formula (1.13) with \( \rho \) replaced by \( \rho + k \), we get the required statement (2.15) under the hypotheses of Theorem 2.8.

**Theorem 2.9.** With the conditions mentioned in (1.18), the following Gegenbauer transform holds:

\[
G^\nu \left[ z^{p-1} W_{\alpha, \beta, \gamma, \delta, \lambda, \sigma}^{(\nu_m \text{me}_N)}(xz; p); n \right] = 2^{2\nu} \left( \frac{2\nu + n - 1}{n} \right) B\left( \nu + \frac{1}{2}, \nu + \frac{1}{2} \right) \sum_{k=0}^{\infty} \mathcal{U}_k \tag{2.18}
\]

where

\[
\mathcal{U}_k = \sum_{k=0}^{\infty} \frac{B_p^{(\nu_m \text{me}_N)}(\nu + \frac{1}{2}, \nu + \frac{1}{2})}{B(\nu + \frac{1}{2}, \nu + \frac{1}{2}) \Gamma(\nu k + \beta) \Gamma(\delta k)} \tag{2.19}
\]

and we assume that the Gegenbauer transform in (2.18) exists.
Remark 2.10. Note that the above transform is precisely the particular case of Theorem 2.8 for \( \alpha = \beta = \nu - \frac{1}{2} \).

For \( \alpha = \beta = 0 \) in Theorem 2.8, we get the Legendre transform as a particular case of Theorem 2.8 given under:

**Theorem 2.11.** With the conditions mentioned with (1.18), the following Legendre transform formula holds true:

\[
\mathcal{L} \left[ z^{\rho-1} W_{a,b,\gamma,\delta,\lambda,\sigma}^{(1)} (xz; p) ; n \right] = 2 \sum_{k=0}^{\infty} \mathcal{U}_k \cdot F_{1:2;1}^{1:1:0} \left[ 1 : -n, n+1; 1 - \rho - k; 1, 1, 1 ; 1, 2 \right] \frac{x^k}{k!} \tag{2.19}
\]

\(|x| < 1; n \in \mathbb{N}_0; \rho \in \mathbb{C} \),

provided that the Legendre transform in (2.19) exists and the coefficients \( \mathcal{U}_k \) are given by (2.16).

3. Derivative Formulas involving generalized Wright function

We recall the Riemann-Liouville fractional derivative of order \( \mu \in \mathbb{C}; \Re(\mu) \geq 0 \) defined by

\[
\mathcal{D}_x^\mu \{ f(x) \} = \frac{1}{\Gamma(r-\mu)} \frac{d^r}{dx^r} \int_0^x f(t) (x-t)^{-\mu+r-1} dt \tag{3.1}
\]

\(( r \in \mathbb{N}; r-1 < \Re(\mu) < r; x > 0 )\).

In particular,

\[
\mathcal{D}_x^\mu \{ f(x) \} = \frac{1}{\Gamma(1-\mu)} \frac{d}{dx} \int_0^x f(t) (x-t)^{-\mu} dt \tag{3.2}
\]

\((0 < \Re(\mu) < 1; x > 0 )\).

A more generalized form of (3.1) introduced by Agarwal et al. \([3]\) is defined as

\[
\mathcal{D}_{x;c,d}^{\mu,\nu} \{ f(x) \} = \frac{1}{\Gamma(r-\mu)} \frac{d^r}{dx^r} \int_0^x f(t) (x-t)^{-\mu+r-1} \mathcal{I}_1 \left[ c, d; \frac{-px^2}{t(x-t)} \right] dt \tag{3.3}
\]

\(( r \in \mathbb{N}; r > \Re(\mu) > r-1; x \in \mathbb{R}^+; p \in \mathbb{R}_0^+; c \in \mathbb{C}; d \in \mathbb{C} \setminus \mathbb{Z}_0^+ \)\).

In particular,

\[
\mathcal{D}_{x;c,d}^{\mu,\nu} \{ f(x) \} = \frac{1}{\Gamma(1-\mu)} \frac{d}{dx} \int_0^x (x-t)^{-\mu} f(t) \mathcal{I}_1 \left[ c, d; \frac{-px^2}{t(x-t)} \right] dt \tag{3.4}
\]

\((p \in \mathbb{R}_0^+ ; 0 < \Re(\mu) < 1; x \in \mathbb{R}^+; c \in \mathbb{C}; d \in \mathbb{C} \setminus \mathbb{Z}_0^+)\).

The special case of (3.1) when and \( p = 0 \) becomes the classical Reimann-Liouville fractional derivative.
The generalization of the extended Reimann-Liouville fractional derivative of a function \( f(x) \) of order \( \mu \) [19] is defined by

\[
\mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{x;\mu,p} f(x) = \frac{1}{\Gamma(-\mu)} \int_0^x f(t) (x-t)^{-\mu-1} \, \mathcal{F} \left[ \{v_m\}_{m \in \mathbb{N}_0}, -\frac{p x^2}{t(x-t)} \right] \, dt \quad (3.5)
\]

When \( \Re(\mu) \geq 0 \), let \( r \in \mathbb{N} \) be the smallest integer greater than \( \Re(\mu) \) and so \( r - 1 \leq \Re(\mu) < r \). Then the generalized Riemann-Liouville fractional derivative of \( f(x) \) of order \( \mu \) is defined by

\[
\mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{x;\mu,p} := \frac{d^r}{dx^r} \mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{x;\mu-p, r-p} f(x)
\]

\[
= \frac{d^r}{dx^r} \left\{ \frac{1}{\Gamma(r-\mu)} \int_0^x f(t) (x-t)^{-\mu+r-1} \, \mathcal{F} \left[ \{v_m\}_{m \in \mathbb{N}_0}, -\frac{p x^2}{t(x-t)} \right] \, dt \right\} \quad (3.6)
\]

We first apply this fractional derivative on a function \( z^\tau \) under the following Lemma:

**Lemma 3.1.** Let \( r - 1 \leq \Re(\mu) < r \) for some \( r \in \mathbb{N} \) and \( \Re(\mu) < \Re(\tau) \). Then we have

\[
\mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{x;\mu,p} \{z^\tau\} = \frac{\Gamma(\tau+1) B_p^{\{v_m\}_{m \in \mathbb{N}_0}}(\tau+1, r-\mu)}{\Gamma(\tau-\mu+1) B(\tau+1, r-\mu)} z^{\tau-\mu}. \quad (3.7)
\]

**Proof.** Applying (3.6) to the function \( z^\tau \), we have

\[
\mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{z;\mu,p} \{z^\tau\} = \frac{d^r}{dz^r} \left\{ \frac{1}{\Gamma(r-\mu)} \int_0^z t^\tau (z-t)^{-\mu+r-1} \, \mathcal{F} \left[ \{v_m\}_{m \in \mathbb{N}_0}, -\frac{p z^2}{t(z-t)} \right] \, dt \right\}.
\]

Setting \( t = z u \) above, we get

\[
\mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{z;\mu,p} \{z^\tau\} = \left( \frac{d^r}{dz^r} z^{r+\tau-\mu} \right) \times \frac{1}{\Gamma(r-\mu)} \int_0^z u^{r+1-1(1-u)-\mu+r-1} \, \mathcal{F} \left[ \{v_m\}_{m \in \mathbb{N}_0}, -\frac{p}{u(1-u)} \right] \, du.
\]

Now, since

\[
\frac{d^r}{dz^r} z^{r+\tau-\mu} = \frac{\Gamma(1+\tau-\mu+r)}{\Gamma(1+\tau-\mu)} z^{\tau-\mu},
\]

therefore, in view of (1.9), we get the desired result. \( \square \)

**Lemma 3.2.** The following easily derivable formulas for the usual beta function \( B(x, y) \) and \( B_p^{c,d}(x, y) \) in (1.8) are mentioned in the lemma below:

\[
B(u, v + 1) = \frac{v}{u + v} B(u, v), \quad (3.8)
\]

\[
B(u + 1, v) = \frac{u}{u + v} B(u, v), \quad (3.9)
\]

\[
B_p^{\{v_m\}_{m \in \mathbb{N}_0}}(u, v) = B_p^{\{v_m\}_{m \in \mathbb{N}_0}}(u + 1, v) + B_p^{\{v_m\}_{m \in \mathbb{N}_0}}(u, v + 1). \quad (3.10)
\]
Now, we present a derivative formula for the generalized Wright function (1.18), asserted by the following theorem:

**Theorem 3.3.** Let \( \Re(\alpha) > 0, p \in \mathbb{R}_0^+, \Re(\delta) > \Re(\gamma) > 0 \) and \( d \in \mathbb{C} \setminus \mathbb{Z}_o^- \). Then

\[
\frac{d^n}{dz^n} W^{(\{v_m\}_m\in\mathbb{N}_0)}_{a,\beta,\gamma,\delta,\lambda,\sigma}(z; p) = \frac{(\gamma)_n (\sigma)_n}{(\delta)_n (\lambda)_n} W^{(\{v_m\}_m\in\mathbb{N}_0)}_{a,na+\beta,\gamma+n,\delta+n,\lambda+n,\sigma+n}(z; p).
\]  

(3.11)

**Proof.** From (1.18), we have

\[
\frac{d}{dz} W^{(\{v_m\}_m\in\mathbb{N}_0)}_{a,\beta,\gamma,\delta,\lambda,\sigma}(z; p) = \sum_{n=1}^{\infty} B_p^{(\{v_m\}_m\in\mathbb{N}_0)} (\gamma + n, \lambda - \gamma) (z)^{n-1} \frac{(\sigma)_n}{(\delta)_n (\lambda)_n}.
\]

Similarly, continuing the process \((n-1)\) times, yields the desired result (3.11).

In terms of (1.18), we obtain

\[
\frac{d}{dz} W^{(\{v_m\}_m\in\mathbb{N}_0)}_{a,\beta,\gamma,\delta,\lambda,\sigma}(z; p) = \frac{\sigma \gamma}{\delta \lambda} \sum_{n=0}^{\infty} (\sigma + 1)_n \frac{B_p^{(\{v_m\}_m\in\mathbb{N}_0)} (\gamma + n + 1, \lambda - \gamma)}{B(\gamma + n + 1, \lambda - \gamma)} \frac{z^n}{n! \Gamma(a + \gamma + \lambda + \beta)}.
\]

(3.12)

Similarly, continuing the process \((n-1)\) times, yields the desired result (3.11).

**Theorem 3.4.** Let \( p \in \mathbb{R}_0^+, x \in \mathbb{R}_0^+ \). Also, let \( r - 2 < \Re(\rho - \delta) < r - 1 \) for \( m \in \mathbb{N} \). Then

\[
\mathfrak{D}_{x;\sigma,d}^{p+1-\delta} \left( x^{p-1} W^{\delta+r-1}_{a,\beta}(x) \right) = \sum_{l=0}^{r} \left( \begin{array}{c} r \\ l \end{array} \right) \frac{\Gamma(\delta + r - 1)}{\Gamma(\delta + l - 1)} \frac{\Gamma(\rho + l)}{\Gamma(\delta + r + l - 1)} \times x^{\delta+l-2} W^{\delta+l+r-1}_{a,\alpha + \beta}(x; p), \quad (r \in \mathbb{N}).
\]

(3.13)

**Proof.** Let \( Q \) denote the left hand side of (3.13). Employing (3.6), we have

\[
Q = \frac{1}{\Gamma(r + \delta - \rho - 1)} \times \frac{d^r}{dx^r} \int_0^x t^{p-1} (x-t)^{\delta+r-p-2} W^{\delta+r-1}_{a,\beta}(t) dt.
\]

Taking \( t = ux \), we obtain

\[
Q = \frac{1}{\Gamma(r + \delta - \rho - 1)}
\]
\[
\times \frac{d^r}{dx^r} \left\{ x^{\delta+r-2} \int_0^1 u^{\rho-1} (1-u)^{\delta+r-\rho-2} W_{\alpha,\beta}^{\delta+r-1} (ux) \mathcal{F} \left[ \{v_m\}_{m \in \mathbb{N}_0}; \frac{-p}{u(1-u)} \right] du \right\}. \tag{3.14}
\]

Applying (2.1) in (3.14), we obtain
\[
Q = \frac{\Gamma(p)}{\Gamma(r + \delta - 1)} \times \frac{d^r}{dx^r} \left\{ x^{\delta+r-2} W_{\alpha,\beta}^{\rho,\delta+r-1;c,d} (x; p) \right\}.
\]

By using the Leibniz's generalized rule for differentiation of product of two functions, we have
\[
Q = \frac{\Gamma(p)}{\Gamma(r + \delta - 1)} \sum_{l=0}^r \binom{r}{l} \frac{d^{(r-l)}}{dx^{(r-l)}} x^{\delta+r-2} \left\{ \frac{d^l}{dx^l} W_{\alpha,\beta}^{\rho,\delta+r-1;c,d} (x; p) \right\}.
\tag{3.15}
\]

Since we have
\[
\frac{d^l}{dx^l} x^a = \frac{\Gamma(a+1)}{a-l+1} x^{a-l} \quad (l \in \mathbb{N}_0)
\tag{3.16}
\]
then, using (3.11) and in view of (3.16), we get the desired result (3.13).

By letting \(r = 1\), we get a particular case of Theorem 3.4 asserted by the following corollary:

**Corollary 3.5.** Let \(p \in \mathbb{R}_0^+, \ x \in \mathbb{R}^+, \ \Re(\rho - \delta) > -1 \) and \(0 < \Re(\rho) < \Re(\delta)\). Then, the undermentioned formula holds:
\[
\mathcal{D}^{\{v_m\}_{m \in \mathbb{N}_0}}_{x;\rho+1-\delta;\rho} \left( x^{\rho-1} W_{\alpha,\beta}^{\delta} (x) \right) = \frac{\Gamma(p)}{\Gamma(\delta)} x^{\delta-2} \left\{ (\delta-1) W_{\alpha,\beta,\rho,\delta}^{\rho,\delta} (x; p) + \frac{\rho x}{\delta} W_{\alpha,\beta+1,\rho,\delta+1}^{\rho,\delta} (x; p) \right\}.
\tag{3.17}
\]

**Theorem 3.6.** Let \(\beta, \gamma, \delta, \sigma \in \mathbb{C}; \ \alpha > -1, \ \Re(\delta) > \Re(\rho) > 0, \Re(\rho) > 0\). Then
\[
\left( \frac{d}{dz} \right)^\kappa \left[ z^{\beta-1} W_{\alpha,\beta,\gamma,\delta,\lambda,\sigma}^{\{v_m\}_{m \in \mathbb{N}_0}} (\mu z^\alpha; p) \right] = z^{\beta-\kappa-1} W_{\alpha,\beta-\kappa,\gamma,\delta,\lambda,\sigma}^{\{v_m\}_{m \in \mathbb{N}_0}} (\mu z^\alpha; p)
\tag{3.18}
\] \((\Re(\beta - \kappa) > 0, \kappa \in \mathbb{N})\).

**Proof.** Using (1.18), and employing term-wise differentiation \(m\) times, we get
\[
\left( \frac{d}{dz} \right)^\kappa \left[ z^{\beta-1} W_{\alpha,\beta,\gamma,\delta,\lambda,\sigma}^{\{v_m\}_{m \in \mathbb{N}_0}} (\mu z^\alpha; p) \right] = \sum_{n=0}^\infty B_p^{\{v_m\}_{m \in \mathbb{N}_0}} (\gamma + n, \lambda - \gamma) (\mu)^n (\sigma)_n (\alpha)_n \left( \frac{d}{dz} \right)^\kappa z^{\alpha_n + \beta - 1}
\]
\[
= \sum_{n=0}^\infty B_p^{\{v_m\}_{m \in \mathbb{N}_0}} (\gamma + n, \lambda - \gamma) (\sigma)_n \mu^n \frac{(\alpha)_n}{n! \Gamma(\alpha + \beta)} \left( \frac{d}{dz} \right)^\kappa z^{\alpha_n - \kappa + \beta - 1}
\]
\[
= z^{\beta-1-\kappa} \sum_{n=0}^\infty B_p^{\{v_m\}_{m \in \mathbb{N}_0}} (\gamma + n, \lambda - \gamma) (\sigma)_n \mu^n \frac{(\alpha)_n}{n! \Gamma(\alpha + \beta - \kappa)} \left( \frac{d}{dz} \right)^\kappa z^{\alpha_n - \kappa + \beta - 1}
\]
\[
= z^{\beta-\kappa} W_{\alpha,\beta-\kappa,\gamma,\delta,\lambda,\sigma}^{\{v_m\}_{m \in \mathbb{N}_0}} (\mu z^\alpha; p),
\tag{3.19}
\]
which is the required result (3.18).

4. Generating relation for the generalized Wright function

In this section, we derive a generating relation involving the product of three generalized Wright functions. We recall the modified form of Exton's result [6] given by

\[ \exp\left(s + t - \frac{zt}{s}\right) = \sum_{u=-\infty}^{\infty} \sum_{v=u^*}^{\infty} s^u t^v H_u^v(z), \]

where,

\[ H_u^v(z) = F_1(-v; u + 1; z) / u! v! = \mathcal{L}_u^v(z) / (u + v)!, \]

and \( \mathcal{L}_u^v(z) \) denotes the Laguerre polynomials (see [10],[11],[20]) and \( u^* = \max(0, -u) \), \( (u = 0, 1, 2, \ldots) \).

**Theorem 4.1.** We derive the following partly bilateral and partly unilateral generating relation for the generalized Wright function

\[ W_{\{m_1 \} m_2 \in \mathbb{N}_0}^{\{m_1 \} m_2 \in \mathbb{N}_0} (s; p) W_{\{m_2 \} m_3 \in \mathbb{N}_0}^{\{m_2 \} m_3 \in \mathbb{N}_0} (t; p) W_{\{m_3 \} m_4 \in \mathbb{N}_0}^{\{m_3 \} m_4 \in \mathbb{N}_0} (z) \]

\[ = \sum_{u=-\infty}^{\infty} \sum_{v=u^*}^{\infty} s^u t^v \mathcal{H}_u^v(z) \]

\[ (i = 1, 2, 3), \quad (4.1) \]

where

\[ \{a, b, \gamma\} H_u^v(z) = \sum_{u=-\infty}^{\infty} \sum_{v=u^*}^{\infty} s^u t^v \times \sum_{k=0}^{\infty} \frac{B_p^{\{m_3 \} m_4 \in \mathbb{N}_0} (\gamma_3 + k, \lambda_3 - \gamma_3) (-z)^k (k_3)_k}{B_3 \lambda_3 - \gamma_3) k! (\beta_3)_{\alpha_3 k} (\delta_3)_{k}} \times \frac{B_p^{\{m_3 \} m_4 \in \mathbb{N}_0} (\gamma_1 + u + k, \lambda_1 - \gamma_1) (\sigma_1)_{u+k}}{B_3 \lambda_1 - \gamma_1) (u + k)! (\beta_1)_{\alpha_1 u+1} (\delta_1)_{u+k}} \times \frac{B_p^{\{m_3 \} m_4 \in \mathbb{N}_0} (\gamma_2 + v - k, \lambda_2 - \gamma_2) (\sigma_2)_{v-k}}{B_3 \gamma_2, \lambda_2 - \gamma_2) (v - k)! (\beta_2)_{\sigma_2} (v - k)! (\delta_2)_{v-k}}. \]

**Proof.** Using the series expansions of (1.18), we write

\[ M(z, s, t) = \sum_{k=0}^{\infty} \frac{B_p^{\{m_3 \} m_4 \in \mathbb{N}_0} (\gamma_3 + k, \lambda_3 - \gamma_3)}{B_3 \lambda_3 - \gamma_3) k! (\alpha_3 k + \beta_3) (\delta_3)_{k}} \times \sum_{i=0}^{\infty} \frac{B_p^{\{m_3 \} m_4 \in \mathbb{N}_0} (\gamma_1 + i, \lambda_1 - \gamma_1)}{B_3 \gamma_1, \lambda_1 - \gamma_1) i! (\alpha_1 i + \beta_1) (\delta_1)_{i}} \times \sum_{j=0}^{\infty} \frac{B_p^{\{m_3 \} m_4 \in \mathbb{N}_0} (\gamma_2 + j, \lambda_2 - \gamma_2) t (j+k) (\sigma_2)_{j}}{B_3 \gamma_2, \lambda_2 - \gamma_2) j! (\alpha_2 j + \beta_2) (\delta_2)_{j}}. \]
Replacing \( i - k \) and \( j + k \) by \( u \) and \( v \) respectively, and rearranging the series that is justified by the absolute convergence, yields

\[
M(z, s, t) = \sum_{u=-\infty}^{\infty} \sum_{v=0}^{\infty} \frac{s^u t^v}{\Gamma(\beta_1)\Gamma(\beta_2)\Gamma(\beta_3)} \times \sum_{k=0}^{\infty} \frac{B_p^{(v+1)}(\gamma_3 + k, \lambda_3 - \gamma_3)(-z)^k(\sigma_3)_k}{B(\gamma_3, \lambda_3 - \gamma_3)k!(\beta_3)_k(\delta_3)_k} \times \frac{B_p^{(v+1)}(\gamma_1 + u + k, \lambda_1 - \gamma_1)(\sigma_1)_{u+k}}{B(\gamma_1, \lambda_1 - \gamma_1)(u+k)!} \times \frac{B_p^{(v+1)}(\gamma_2 + v - k, \lambda_2 - \gamma_2)(\sigma_2)_{v-k}}{B(\gamma_2, \lambda_2 - \gamma_2)(v-k)!}.
\]

\( \square \)

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