Diffuse Gamma Rays in 3D Galactic Cosmic-ray Propagation Models

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Abstract. The \textsc{Picard} code for the numerical solution of the Galactic cosmic ray propagation problem allows for high-resolution models that acknowledge the 3D structure of our Galaxy. \textsc{Picard} was used to determine diffuse gamma-ray emission of the Galaxy over the energy range from 100\,MeV to 100\,TeV. We discuss the impact of a cosmic-ray source distribution aligned with the Galactic spiral arms for a range of such spiral-arm models. As expected, the impact on the gamma-ray emission is most distinct in the inverse-Compton channel, where imprints of the spiral arms are visible and yield predictions that are no longer symmetric to the rotational axis of the Milkyway. We will illustrate these differences by a direct comparison to results from previous axially symmetric Galactic propagation models: we find differences in the gamma-ray flux both on global scales and on local scales related to the spiral arm tangents. We compare gamma-ray flux and spectra at on-arm vs. off-arm projections and characterize the differences to axially symmetric models.

INTRODUCTION

The fact that our Galaxy does not have an azimuthally symmetric shape is lately being acknowledged in recent Galactic cosmic-ray propagation models [see, e.g., 1, 2, 3, 4, 5, 6]. These models focus on the observation that our Galaxy has a spiral structure [see 7] and implications for the transport of cosmic rays within the Galaxy. This respective spiral structure will manifest itself in the distribution of ISM gas [see 8] and the interstellar magnetic field [see, e.g. 9, 10] but also in the distribution of cosmic-ray sources [see, e.g. 1, 2, 11].

The impact of the source distribution on the cosmic-ray flux and distribution within the Galaxy has been studied by several groups [see 1, 2, 4, 12, 5, 13, 6]. It was shown that a good agreement with the observed cosmic-ray data at Earth [see 1, 13] can be found both in models with an axisymmetric source distribution and in models using a source distribution motivated by the observed three-dimensional structure of our Galaxy. Different spatial distributions of the sources lead to different spatial distributions of the cosmic-ray flux, which should lead to distinct effects on the emission of radiation by the interaction of cosmic rays with the interstellar medium resulting in the different models.

Here, we focus on the imprint of non-axisymmetric source distribution on the gamma-ray emission from our Galaxy. For this we use simulations performed with the \textsc{Picard} code. This code was introduced in [14] and is optimised for such spatially three-dimensional numerical modelling. In the following we introduce the \textsc{Picard} code with a focus on recent enhancements. Subsequently we discuss the ensuing gamma-ray emission for our models and show the related energy spectra.

THE PICARD CODE

The \textsc{Picard} code was introduced in [14] with a focus on spatially three-dimensional numerical modelling of Galactic cosmic-ray transport problems. The code is optimised to find steady-state solutions to the cosmic-ray transport
equation:
\[
\frac{\partial \psi_i}{\partial t} = q(\vec{r}, p) + \nabla \cdot \mathbf{D} \nabla \psi_i + \frac{\partial}{\partial p} p^2 D_{pp} \frac{1}{p} \frac{\partial}{\partial p} \psi_i - \nabla \cdot (\vec{v} \psi_i) - \frac{\partial}{\partial p} \left( \dot{p} \psi_i - \frac{2}{3} \left( \nabla \cdot \vec{v} \right) \psi_i \right) - \frac{1}{\tau_f} \psi_i - \frac{1}{\tau_r} \psi_i, \tag{1}
\]
but it can also be applied to solve time-dependent problems. A numerical solution to the steady-state problem is found by discretising the differential operators in equation (1) in space and momentum and then finding an approximate solution to the resulting linear system of equations. In P\textsc{icard} the linear system of equations can be solved using different standard techniques, including a range of multigrid methods [see 15] or the BiCGStab method as described in [16], with the possibility to add other solvers into the modular framework of the code.

P\textsc{icard} features all relevant transport processes inherent in equation (1): spatial diffusion, spatial convection, momentum diffusion and relevant energy-loss processes. Fluxes of light species depend on those of heavier species due to the possible fragmentation or radiative-decay processes. This is realised through the fragmentation and radiative-decay timescales $\tau_f$ and $\tau_r$, respectively, together with a consideration of a related source term for other species. The resulting nuclear network is efficiently implemented in P\textsc{icard}, leading to a numerically optimised and accurate treatment of coupled nuclear species [see 13].

Prominent features of the P\textsc{icard} code include the possibility to consider anisotropic spatial diffusion [for a discussion see, e.g. 17, 18] and arbitrary convection velocities $\vec{v}$. The numerical solver is MPI-parallel, thus facilitating the possibility to achieve up to deca-parsec scale resolution by making use of large-scale distributed-memory computers. The numerical solution computed using P\textsc{icard} is acquired with a user-defined accuracy without the need to check the quality of the numerical solution a-posteriory. P\textsc{icard} is continuously expanded where, e.g., the currently ongoing implementation of an improved interstellar radiation field is discussed by [19].

With this code we studied the implications of spiral-arm cosmic-ray source distributions on the spatial variability of the cosmic-ray flux in [5]. Additionally, we showed in [13] that the transport parameters in equation (1) can be adapted to such localised source distributions, thus leading to a good agreement between predicted cosmic-ray fluxes at Earth and corresponding observations. For the secondary to primary ratios there appears a dependence on the distance to the nearest sources in modules using such localised sources. In the following we will address the impact of these models on the Galactic diffuse gamma-ray emission.

**GAMMA-RAY EMISSION**

In [13] we found that a Galactic cosmic-ray propagation model featuring localised sources can reproduce cosmic-ray measurements at Earth as well as a model with underlying axisymmetry. In contrast to such an axisymmetric model, a spiral-arm source distribution leads to a localisation of the cosmic-ray flux, at least for the primary cosmic-ray species. The distinct difference in spatial distribution between the different models can be expected to lead to differences in the resulting gamma-ray flux in the different models. In the following we focus on the comparison between a model using an axisymmetric source distribution and one using a spiral-arm source distribution with four spiral arms. Details can be found in [5] and [13], where this specific model is referred to as the Steiman-model.

The relevant channels for diffuse gamma-ray emission are $p^0$-decay, bremsstrahlung, and inverse Compton (IC) radiation. The emissivity of the former two of these processes results from the product of the gas-density with the respective cosmic-ray density – in this case nucleons and electrons, respectively. Therefore, it is not easy to distinguish whether a spatial feature in the gamma-ray emission results from local structures in the gas distribution or the cosmic ray distribution. Since current models of the Galactic radiation field are spatially much smoother than the gas distribution [see 20, 21, 22], the IC channel is a much clearer tracer of the localisation of the sources within our models. This is obvious when investigating the gamma-ray emission in Fig. 1. A comparison shows that the Galactic plane near the Galactic centre becomes slightly brighter for the model with a spiral-arm source distribution. Otherwise localised features affected by the different source models are not readily apparent.

In the following we will focus on the region near the Galactic centre, where the localised features of the cosmic-ray distribution in the Steiman-model have the most obvious impact on the gamma-ray emission. This is caused by the spiral arms that are tightly wound around the Galactic centre leading to several spiral arm tangents near the Galactic centre along which the projected source intensity is enhanced. This is shown in Fig. 2 where the projected source intensity near the Galactic centre region is shown. The specific spiral-arm tangents are consistent with observations [7] since the Steiman model is driven by observations of the Galactic spiral arms [23]. The specific spiral-arm tangents visible in Fig. 2 are tangents to the Sagittarius / Carina arm at $l \sim -50^\circ$, $l \sim -16^\circ$, and $l \sim -75^\circ$, to the Scutum / Crux Centaurus arm at $l \sim 30^\circ$ and $l \sim -50^\circ$, the Norma arm at $l \sim 19^\circ$ and at $l \sim -30^\circ$, and the Perseus arm at $l \sim -22^\circ$. 
FIGURE 1: All-sky gamma-ray emission at $E_\gamma = 10$ TeV for an axisymmetric (top) and a four-arm (bottom) source distribution. The gamma-ray emission is shown in units of $\text{cm}^{-2} \text{s}^{-1} \text{sr}^{-1} \text{MeV}^{-1}$.

In Fig. 2 we also show the IC emission at an energy of $\sim 100$ GeV. The increase in gamma-ray flux at the positions of the spiral-arm tangents is obvious at this energy. High-energy electrons are tightly coupled to their sources due to their very high energy losses [5]. Therefore, the related gamma-ray emission traces these source regions leading to an imprint of the spiral-arm tangents also on the gamma-ray emission. In the total gamma-ray emission these effects are not so obvious, because of the impact of the gas-distribution discussed above and also because the Galactic diffuse gamma-ray emission is in many directions dominated by $\pi^0$-decay emission. In the following we investigate the quantitative impact of the spiral-arm source distribution on the gamma-ray emission.

GAMMA-RAY SPECTRA

For a quantitative analysis we compare gamma-ray spectra at on-arm and off-arm positions. In Fig. 3 corresponding results are shown for spiral-arm tangents of the Norma and the Sagittarius arm. Obviously, the on-arm flux is significantly higher in both cases and for all gamma-ray emission channels. Especially, for Bremsstrahlung and IC the on-arm spectra are also harder. For these channels we found that the power-law index can increase by 0.1 or even more at a spiral-arm tangent as compared to an off-arm position. This is somewhat more pronounced for the Sagittarius arm, because in the case of the Norma arm the off-arm region is in a region featuring several other spiral-arms. Thus, the Norma off-arm region is contaminated by the presence of other close-by cosmic-ray sources.

The spectral hardening of the IC flux can be explained by the energy dependence of the electron’s energy losses. Due to their rapid increase with energy, high-energy electrons suffer from the most severe energy losses. Therefore, high-energy electrons can only be found in the close vicinity of their sources. In [5] a significant hardening of the electron spectrum was observed for on-arm positions. This naturally translates into a hardening of the resulting photon spectrum for a sight line in the direction of a spiral-arm tangent.

The $\pi^0$ channel shows only mild change in spectral slope. This is also consistent with the observations in [5] where the proton flux was found to mainly change in amplitude from an on-arm to an off-arm position. This reflects the fact that the spatial distribution of nuclear cosmic rays is mainly driven by spatial diffusion due to their rather low energy losses. Also in gamma rays this channel mainly features a change in amplitude when comparing an on-arm...
FIGURE 2: Left: Projected source intensity for the four-arm source-distribution model. Spiral-arm tangents are visible as enhanced intensity in the Galactic plane. Right: Inverse Compton gamma-ray emission at 100 GeV for the same model.

FIGURE 3: Spectra for the different gamma-ray emission channels for two of the spiral-arm tangents. In each plot on-arm and off-arm data are compared, where the on-arm spectra are the ones with the higher fluxes. On the left results are shown for a tangent of the Norma arm with the on-arm data taken at -30° longitude and the off-arm data taken at -35° longitude. On the right corresponding data for a Sagittarius-arm tangent are shown with the on-arm data taken at 50° longitude and the off-arm data taken at 62° longitude. Results are integrated over a circle with a radius of 2°.

and an off-arm region.

Since the total flux in both cases is dominated by π⁰ decay, the change in slope for the total flux is also rather small. Only in the energy regime from a few up to a few hundred GeV, where the contribution of IC emission to the total gamma-ray flux is largest, the spectrum at an on-arm position becomes harder with a change in power law index of up to 0.05. In the present case we used the same source-distribution for electrons and nuclei. In a model using different source distributions the corresponding changes would become larger, for such localised source distributions as were discussed in this case. In fact, for high-energy electrons it will be important to consider the actual sources [see 24, 25, 26] when aiming for an accurate model for the diffuse gamma-ray emission in the Galactic plane. Our model shows that localised sources will not only have an impact on the gamma-ray emission in the Galactic plane, but can also affect the local spectral slope. This will make a template-fitting approach at least more problematic than a physically motivated modelling approach based on the propagation physics of Galactic cosmic rays. One should be aware that in the current approach only the distribution of the cosmic-ray sources was adapted to the Galactic structure. With an adoption of the gas-, magnetic field-, and possibility the interstellar radiation distribution the differences between an axisymmetric approach and one acknowledging the actual structure can become even larger, holding the possibility to also explain the small-scale structure in the gamma-ray emission in the Galactic plane.
CONCLUSIONS

In this study we showed that Picard is ready to predict gamma-ray fluxes for a comparison with data in the context of Galactic cosmic-ray transport models taking into account the three-dimensional structure of our Galaxy. We showed a first quantitative discussion of the impact of localised sources on the Galactic diffuse gamma-ray emission. We found a hardening of the gamma-ray emission spectrum in the direction of the cosmic-ray sources for the leptonic emission channels. Since in our model the emission is dominated by π0 decay the change in slope for the total gamma-ray flux is rather small and also limited to energies between 1 and a few hundred GeV, where the leptonic channels provide a comparatively large fraction of the total emission. The current analysis was focussed on the impact of the source distribution. While this can only be viewed as a first step, with the future need to include the full Galactic structure for gas-, magnetic field-, and interstellar radiation-distribution, only a dissemination of the individual effects can show the respective impact on changes in the gamma-ray emission.

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