Identification of Ne\textsuperscript{VIII} lines in H-deficient (pre-) white dwarfs: a new tool to constrain the temperature of the hottest stars\textsuperscript{*} **

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ABSTRACT

For the first time, we have identified Ne\textsuperscript{vii} absorption lines in far-UV spectra of the hottest known (\(T_{\text{eff}} \gtrsim 150 000\) K) hydrogen-deficient (pre-) white dwarfs of spectral type PG1159. They are of photospheric origin and can be matched by synthetic non-LTE line profiles. We also show that a number of UV and optical emission lines in these stars can be explained as being photospheric Ne\textsuperscript{vii} features and not, as hitherto suspected, as ultrahigh ionized O\textsuperscript{vii} lines created along shock-zones in the stellar wind. Consequently, we argue that the long-standing identification of the same emission lines in hot [WR]-type central stars as being due to ultrahigh-ionised species (O\textsuperscript{vii-viii}, C\textsuperscript{v-vii}) must be revised. These lines can be entirely attributed to thermally excited species (Ne\textsuperscript{vii-viii}, N\textsuperscript{v}, O\textsuperscript{vii}). Photospheric Ne\textsuperscript{vii} lines are also identified in the hottest known He-rich white dwarf (KPD\,0005+5106), some of which were also attributed to O\textsuperscript{vii} previously. This is a surprise because it must be concluded that KPD\,0005+5106 is much hotter (\(T_{\text{eff}} \approx 200 000\) K) than hitherto assumed (\(T_{\text{eff}} \approx 120 000\) K). This is confirmed by a re-assessment of the H\textsc{ii} line spectrum. We speculate that the temperature is high enough to explain the mysterious, hard X-ray emission (1 keV) as being of photospheric origin.

Key words. Stars: abundances – Stars: atmospheres – Stars: evolution – Stars: AGB and post-AGB – White dwarfs

1. Introduction

Observations of extremely hot post-AGB stars (\(T_{\text{eff}} \gtrsim 100 000\) K) with the Far Ultraviolet Spectroscopic Explorer (FUSE) have led to the discovery of many unidentified spectral lines. Their identification is difficult, mostly because of the lack of accurate atomic data for highly ionized elements. On the other hand, their identification is rewarding because they are often the only features by which particular species are accessible for abundance determinations. Such species can be used to test stellar evolutionary models. For example, several lines of fluorine (F\textsuperscript{v-vi}), neon (Ne\textsuperscript{v}), silicon (Si\textsuperscript{v}) and argon (Ar\textsuperscript{vii}) were identified for the first time (Jahn et al. 2007; Werner et al. 2004a, 2005, 2007). In particular, the Ne\textsuperscript{v} \(\lambda\) 973 Å line occasionally displays a prominent P Cygni profile, which is an important tool for stellar wind analyses (Herald et al. 2005). In this paper, we report the identification of a number of Ne\textsuperscript{v} lines in the FUSE spectra of the hottest known hydrogen-deficient (pre-) white dwarfs. We show that the mere occurrence of such highly ionized neon lines puts a strict lower limit on the effective temperature (\(\approx 150\,000\) K).

Ne\textsuperscript{vii} is a one-valence electron, lithium-like ion with a relatively simple energy-term structure, however, a detailed compilation of accurate level energies and line wavelengths, which is a prerequisite of any quantitative work, became available only recently (Kramida & Buchet-Poulizac 2006). Our model atmosphere and spectrum synthesis calculations predict the presence of Ne\textsuperscript{vii} lines in other wavelength regions, too. As a surprising result, we find that all previous identifications of ultra-high ionization (i.e., non-thermally excited) oxygen lines (O\textsuperscript{vii}) in UV and optical spectra of PG1159 stars, and the hottest known DO white dwarf, are wrong. We find that, instead, these lines are due to photospheric Ne\textsuperscript{vii}. The earlier identification of these lines as O\textsuperscript{vii} was motivated by the supposed occurrence of the same lines in early-type Wolf-Rayet central stars (i.e., spectral type [WCE]) and it was argued that they are formed in shocked wind regions. The results presented here provide evidence that also in [WCE] stars these, as well as other lines that were assigned to ultrahigh-ionized C and O, do probably stem from thermally excited neon.

We present our observations and line identifications in Sect.\textsuperscript{2} and describe our modeling in Sect.\textsuperscript{3}. The results from line profile fits to individual objects are presented in Sect.\textsuperscript{4}. Implications for [WCE] stars are discussed in Sect.\textsuperscript{5} and we conclude with Sect.\textsuperscript{6}.

2. Observations and Ne\textsuperscript{VIII} line identifications

FUSE observations and data reduction for most of our program stars were described in our previous work (Werner et al. 2004a, 2004b, 2005, 2007). Table\textsuperscript{1} lists our program stars with references to results from previous analyses. The FUSE instrument consists of four coaligned telescopes, each with a prime-focus
spectrograph. Descriptions of the FUSE instrument, and channel alignment and wavelength calibration issues, are given by Moos et al. (2000) and Sahnow et al. (2000). The FUSE spectra cover the wavelength range from the Lyman edge up to 1187 Å with a spectral resolution of 0.05 Å. For spectra that are too faint to permit co-alignment of individual exposures, the spectral resolution may be degraded to 0.1 Å. The S/N ratio in the LiF2a spectra in the vicinity of the Ne viii lines being studied here ranges from a low of 17:1 per 0.05 Å resolution element for Longmore 4 to 70:1 per resolution element for RX J2117.1+3412; it was greater than 45:1 for the other objects. The S/N ratio in the LiF1b spectra at these wavelengths was typically ~30% lower. UV spectra of KPD 0005+5106 and RX J2117.1+3412, taken with the FOS and GHRS spectrographs aboard the Hubble Space Telescope, were retrieved from the MAST archive. High-resolution optical spectra of K1-16, RX J2117.1+3412, KPD 0005+5106, and H1504+65 were obtained at the Keck observatory and the HIRES spectrograph. For details on data reduction see Zuckerman & Reid (1998). Spectra of NGC 246 and Longmore 4 were obtained with the ESO Very Large Telescope and the UVES spectrograph, in the framework of the SPY project (Napiwotzki et al. 2003).

The starting point of our analysis was the identification of two strong and broad absorption features of Ne viii at \( \lambda = 1162 - 1165 \) Å in the FUSE spectra of five PG1159 stars and the DO white dwarf KPD 0005+5106 (left panel of Fig. 1). They are blends of several \( n = 5 \rightarrow 6 \) lines with high angular quantum number \( l \). Ne viii lines are not detected in the FUSE spec-

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**Fig. 1.** Left panel: Identification of Ne viii lines in the FUSE spectra of PG1159 stars and the DO white dwarf KPD 0005+5106. Several lines of the \( n = 5 \rightarrow 6 \) transition are detected, as labeled in detail at the uppermost spectrum. Overplotted are computed profiles with \( T_{\text{eff}} \) and \( \log g \), as indicated. Middle and right panels: Optical spectral regions where we identified the Ne viii \( n = 8 \rightarrow 9 \) and \( n = 9 \rightarrow 10 \) transitions. All observed spectra are shifted such that the photospheric lines appear at their rest wavelengths.
tra of any other PG1159 star, indicating that the minimum $T_{\text{eff}}$ for exhibiting these lines is around 150 000 K (this limit will be assessed more closely below). In addition to these lines, several weaker Ne viii lines are detectable in the FUSE spectra of some objects, particularly in RX J2117.1+3412 (panels a)–c) in Fig. 2. These are low-$l$ $n = 5 \to 6$ lines and an isolated $n = 6 \to 8$ line. The strongest, high-$l$, $n = 6 \to 7$ transitions are located at $\lambda$ 1932 Å. We found only one archival spectrum covering this wavelength position, namely an HST/FOS observation of KPD 0005+5106. An emission feature is barely detectable at this location (panel d) of Fig. 2), formerly attributed to O viii. Panel e) displays HST/GHRS spectra with a Ne viii $n = 7 \to 8$ emission line that was also thought to stem from O viii. The adjacent emission feature in KPD 0005+5106 is from nitrogen; it is not included in the model.

Table 1. Effective temperature and surface gravity of the program stars as taken from the literature. As described in the text, the discovery of Ne viii lines suggests that $T_{\text{eff}}$ for KPD 0005+5106 and Longmore 4 must be significantly higher.

| Object       | Spectral Type | $T_{\text{eff}}$ [K] | log $g$ (cgs) | Reference      |
|--------------|---------------|-----------------------|---------------|----------------|
| H1504+65     | PG1159        | 200 000               | 8.0           | A              |
| RX J2117.1+3412 | PG1159     | 170 000               | 6.0           | B              |
| NGC 246      | PG1159        | 150 000               | 6.7           | C              |
| K1-16        | PG1159        | 140 000               | 6.5           | D              |
| NGC 2371     | [WCE]         | 135 000               | 5.5           | E              |
| Longmore 4   | PG1159        | 120 000               | 5.5           | F              |
| KPD 0005+5106 | DO           | 120 000               | 7.0           | G              |

References: A: Werner et al. 2004b, B: Werner et al. 1996b, C: Rauch & Werner 1997, D: Werner et al. 1992, E: Herald & Bianchi 2004, F: Rauch & Werner 1997, G: Werner et al. 1994

Table 2. Ne viii lines identified in our program stars. Wavelengths are computed from NIST level energies (except for the $5p \to 6d$, $5f \to 6d$, and $6p \to 8d$ transitions, see text) and given in vacuum or air for $\lambda$ smaller or larger than 3000 Å, respectively. The presence of the feature at $\lambda$ 1932 Å is uncertain.

| Wavelength / Å | Transition |
|----------------|------------|
| 1060.36, 1061.03 | $5s \to 6p$ |
| 1129.02, 1130.12, 1130.35 | $5p \to 6d$ |
| 1150.12, 1150.81, 1150.92 | $6p \to 8d$ |
| 1162.24, 1162.67 | $5d \to 6f$ |
| 1164.54, 1164.75 | $5f \to 6g$ |
| 1165.88 | $5g \to 6h$ |
| 1165.94, 1166.15, 1166.18 | $5f \to 6d$ |
| 1170.04, 1170.29 | $6d \to 8f$ |
| (1932.04) | $6h \to 7i$ etc. |
| 2976.75 | $7i \to 8k$ etc. |
| 4340.77 | $8k \to 9l$ etc. |
| 6068.63 | $9l \to 10m$ etc. |

3. Model atmospheres and neon line formation calculations

We use a grid of line-blanketed non-LTE model atmospheres, which is described in detail in Werner et al. (2004a). In essence, the models include the main photospheric constituents, namely, $\ldots$

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1. [http://physics.nist.gov/PhysRefData/ASD/index.html](http://physics.nist.gov/PhysRefData/ASD/index.html)
He, C, O, and Ne. The neon model atom consists of the ionization stages Ne IV–IX. Of particular importance for the investigations presented in this paper is the Ne VII ion. It consists of 5 NLTE levels and 6 lines, which is sufficient for the atmospheric model structure computation, but considerable extensions were required in order to be able to compute profiles of the identified highly excited lines. Since we aimed at the computation of lines involving the \( n = 10 \) levels (the \( n = 9 \rightarrow 10 \) emission at 6068 Å), we included all 77 levels up to \( n = 12 \) with all 511 line transitions between them. In addition, levels up to \( n = 14 \) are included as LTE levels. With this extended Ne VII model ion, we performed line-formation iterations for neon, i.e., keeping the atmospheric structure fixed. For the final spectrum synthesis, fine-structure splitting of levels and lines was considered, if possible. Energies for fine-structure splitting are not available for all levels. In particular, they are only partly available for levels involved in the \( n = 5 \rightarrow 6 \) absorption lines at \( \lambda 1162 – 1165 \) Å. As a consequence, the detailed structure of the computed line profile cores within these absorption troughs cannot exactly match the observations and, indeed, a close inspection of Fig. 1 reveals this problem. Level energies were taken from NIST (Kramida & Buchet-Poulizac 2006). Oscillator strengths were taken from the Opacity (Seaton et al. 1994) and IRON (Hummer et al. 1993) Projects databases (TIOPTOBase2), which are, however, not complete for all transitions of our model atom. For lines involving levels with \( n = 11 \) and 12, we extrapolated oscillator strengths from transitions to lower quantum numbers. For a lithium-like ion this appears reasonably accurate; the estimates are probably within a 10% error. Photoionization cross-sections are taken from theOpacity Project database when available or, otherwise, computed in hydrogen-like approximation. Electron collisional rates were calculated with usual approximation formulae. The Ne model atoms that were used for this analysis have been developed in the framework of the German Astrophysical Virtual Observatory (GAVO3) project and are provided within the Tübingen Model-Atom Database TMAD4.

The treatment of Stark broadening of Ne VII lines poses a severe problem. We employ an approximate formula to compute Stark widths, which we use routinely to compute profiles of lines between highly excited levels of lithium-like ions (C IV, O v, Werner et al. 1991), and which accounts for linear Stark broadening. While this formula gave reasonable results for the other Li-isoelectronic ions in our past work, we found that the \( n = 5 \rightarrow 6 \) absorption lines in the FUSE spectral range came out much too narrow (by about a factor of three) compared to all of the observations. Interestingly, a similar notorious problem is encountered in the interpretation of Ne VII line widths of laboratory spectra. Measurements of the \( 3s \rightarrow 3p \) transition at 2821 Å (Glenzer et al. 1992) display line profiles that are roughly a factor of two wider than predicted by theory, i.e., the \( Z^{−2} \) scaling of the line width, as predicted by the impact theory, significantly deviates from observations. Even today, this discrepancy is not resolved, despite of intense works with both improved experiments and quantum-mechanical calculations (Hegazy et al. 2003; Griem & Raclchenko 2006). Our finding from stellar spectra indirectly suggests that the line widths measured in laboratory are in fact correct, and that instead, the line-broadening theory is still inaccurate. It is necessary to note, however, that we observe different lines of Ne VIII.

To cope with this problem we formally reduced the ionic core charge of Ne VII in our line broadening formula from \( Z = 8 \) to \( Z = 2 \). This yields widths of the \( n = 5 \rightarrow 6 \) lines that coincide with the observations. Of course, this procedure is rather unsatisfactory and would prohibit an accurate abundance determination from these lines. In addition, a close inspection of the computed profiles for the optical lines (i.e., the higher-\( n \) transitions, Fig. 1) reveals that they become too broad with our crude Z-reduction. Fortunately, the neon abundance of all PG1159 stars discussed here is known to be of the order 2% (Werner et al. 2004a). The focus of our paper is on the extreme temperature sensitivity of the Ne VII lines and line broadening is a subordinate effect only.

4. Results

We now compare our computed Ne VII line profiles to the observations (Fig. 1). It is instructive to look at Fig. 1 in which we demonstrate that the strengths of the \( \lambda \lambda 1162 – 1165 \) Å absorption lines strongly depend on \( T_{\text{eff}} \) and \( \log g \). The minimum \( T_{\text{eff}} \) for a line detection is a strong function of gravity, ranging from about 160 000 K at \( \log g = 6.5 \) to 140 000 K at \( \log g = 5.7 \). The mere detection of these Ne VII lines thus provides a strict lower limit to \( T_{\text{eff}} \). Generally, the overall characteristics of the Ne VII lines are reproduced by the models. We find that the UV lines are in absorption, while the optical lines are in emission.

Let us discuss in some detail each of the six objects in which we detected the UV Ne VII lines, particularly those.

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2 http://vizier.u-strasbg.fr/topbase/
3 http://www.g-vo.org
4 http://astro.uni-tuebingen.de/~rauch/TMAD/TMAD.html

Fig. 3. The Ne VII lines are strongly sensitive to \( T_{\text{eff}} \) and \( \log g \). Generally, they become stronger with increasing \( T_{\text{eff}} \) and decreasing \( \log g \). Top panel: \( T_{\text{eff}} \)-sensitivity in DO models with \( \log g = 6.5 \). The lines disappear if \( T_{\text{eff}} \) is below 160 000 K. Middle panel: \( T_{\text{eff}} \)-sensitivity in PG1159 models with lower gravity (\( \log g = 5.7 \)). In this case, the lines are detectable down to \( T_{\text{eff}} = 140 000 \) K. Bottom panel: \( \log g \) sensitivity in PG1159 models at \( T_{\text{eff}} = 160 000 \) K. The gravity must be sufficiently low, otherwise the Ne VII lines are not detectable. The numbers in the lower right corner of the panels give the model abundances of He/C/O/Ne in % mass fraction.
Fig. 4. The He II lines in KPD 0005+5106 are compared with two different models. Left panel: $T_{\text{eff}}=200\ 000$ K, $\log g=6.5$. Right panel: $T_{\text{eff}}=120\ 000$ K, $\log g=7.0$. The hotter model yields a much better fit. In particular, only this model matches the emission cores in the $\lambda\lambda$ 6560 and 4686 Å lines. The emission line at $\lambda$ 4340 Å is due to Ne VIII and is only exhibited in the hotter model.

4.1. PG1159 stars

Models for individual objects were computed with element abundances as given in Werner et al. (2005). As already mentioned, the neon abundance is kept fixed at 2%. We started with values for $T_{\text{eff}}$ and $\log g$ also taken from this reference; typical uncertainties are 10% and 0.5 dex, respectively.

K1-16 The previously determined parameters are $T_{\text{eff}}=140\ 000$ K and $\log g=6.4$. Clearly, a model with these values shows no Ne VIII lines at all. We find an acceptable fit only at $T_{\text{eff}}$ and $\log g$ values that are slightly beyond the error bars, namely $T_{\text{eff}}=160\ 000$ K and $\log g=5.8$.

RX J2117.1+3412 We find a good fit to the Ne VIII lines with models close to the literature values ($T_{\text{eff}}=170\ 000$ K, $\log g=6.0$). At these parameter values, the Ne VIII absorption lines are strongest. This explains why we can identify a number of additional weak Ne VIII lines in the FUSE spectrum (Fig. 2).

NGC 246 A good fit to the Ne VIII lines is obtained at the literature values ($T_{\text{eff}}=150\ 000$ K, $\log g=5.7$). Note that NGC 246 is a fast rotator and the computed profiles were broadened with 70 km/s (Rauch & Werner 1997).

Longmore 4 The strong Ne VIII lines in the FUSE spectrum suggest that this star is as hot as RX J2117.1+3412. We find a good fit with a $T_{\text{eff}}=170\ 000$ K, $\log g=6.0$ model. Therefore, Longmore 4 is significantly hotter than previously thought ($T_{\text{eff}}=120\ 000$ K, $\log g=5.5$). This result is, however, not surprising because we have already found independent hints that the temperature is at least about 150 000 K (Werner et al. 2004a). Note that Longmore 4 exhibited much more prominent optical

cases in which contradictions to $T_{\text{eff}}$ determinations from previous work appear. This mainly affects the DO white dwarf KPD 0005+5106, whereas for the PG1159 stars the Ne VIII line features can be matched with model parameters that are in good or acceptable agreement with previous results.
Ne v ly lines during its observed “outburst” and change of spectral type from PG1159 to [WCE] (Werner et al. 1992).

H1504+65  The Ne v ly lines in the FUSE spectrum are shallow and broad. This is a consequence of the high gravity (8.0). The profiles confirm the extremely high \( T_{\text{eff}} \) (200 000 K).

4.2. The DO white dwarf KPD 0005+5106

An analysis of optical and HST UV spectra of KPD 0005+5106 gave the result \( T_{\text{eff}} = 120 000 \) K and \( \log g = 7 \) (Werner et al. 1994). Although being the hottest known DO white dwarf, the identification of Ne v ly lines is a big surprise. Figure 2 shows that at this relatively high gravity the temperature must be around 200 000 K. We find that models with \( T_{\text{eff}} = 180 000 \) to 200 000 K and \( \log g = 6.5 \) to 7.0 give the best fit to the Ne v ly lines. Interestingly, this is not in contradiction with the He ii lines. We even find that a model with \( T_{\text{eff}} = 200 000 \) K and \( \log g = 6.5 \) gives better He ii line fits than the cooler 120 000 K model. This is shown in Fig. 2. To be more specific, the hot model is able to match the height of the central emissions reversals in the \( \lambda \lambda 4686 \) and 5650  \( \AA \) line cores. At the same time, the \( \lambda 1640 \)  \( \AA \) line is matched very well. We conclude that KPD 0005+5106 is significantly hotter than previously thought.

Unlike for the PG1159 stars, the neon abundance in KPD 0005+5106 is not known from previous analyses. The usual diagnostic lines are Ne v \( \lambda \lambda 973 \) and 3644  \( \AA \). The FUSE spectrum of KPD 0005+5106 is contaminated by interstellar H\(_2\) absorption, making the detection of Ne v \( \lambda 973 \)  \( \AA \) impossible. In a medium-resolution optical spectrum published by Werner et al. (1994), no line feature is detected at \( \lambda 3644 \)  \( \AA \). We have verified that this is compatible with all our models presented here in the \( T_{\text{eff}} = 120 000 \) to 200 000 K range.

We assumed that the Ne abundance is determined by radiative levitation. The hottest model presented by Chayer et al. (1995) has \( T_{\text{eff}} = 100 000 \) K with \( \log g = 7.5 \) and predicts \( \log (\text{Ne}/\text{He}) = -4 \) by number. It is difficult to extrapolate this to a significantly higher \( T_{\text{eff}} \) but the tendency is that the Ne abundance increases with \( T_{\text{eff}} \). A further increase can be expected because of the smaller gravity. Although uncertain, it is not unreasonable that we set Ne = 1% (by mass), that is, \( \log (\text{Ne}/\text{He}) = -2.7 \) by number. We stress that the high \( T_{\text{eff}} \) derived from the Ne v lines is hardly affected by this assumption. In addition, the He ii line profiles are essentially independent of the Ne abundance.

5. Implications for ultrahigh-ionisation emission line identifications in [WCE] stars

Since the supposed identification of several optical emission lines in the hot [WC] star Sand 3 as being due to ultrahigh ionised (i.e. non-thermally excited) C and O (C v and O v-iv) (Barlow et al. 1980), a number of authors claimed the identification of these features in other early-type [WC] stars. In the light of our results, we propose that most, and probably all, of these features stem from thermally excited ionisation stages of neon (Ne v-iv), carbon and oxygen. Table 3 summarizes our proposed identifications. We give qualitative arguments for this, but detailed NLTE modeling with expanding model atmospheres will be necessary for a quantitative confirmation.

A broad emission feature in Sand 3 and other [WCEs] at \( \lambda 3893 \)  \( \AA \) was attributed to \( n = 7 \rightarrow 8 \) transition of O v. We propose that it originates from a prominent Ne v feature recently identified in the hottest PG1159 central stars (Werner et al. 2004a), which has its strongest component located at \( \lambda 3892 \)  \( \AA \). In addition, the Ne v \( n = 7 \rightarrow 8 \) transition contributes. The emission at \( \lambda 4340 \)  \( \AA \) is probably from Ne v, not O v. Emission lines at \( \lambda 4555 \)  \( \AA \) and 5665  \( \AA \) were assigned to the O v transitions \( n = 9 \rightarrow 11 \) and \( n = 8 \rightarrow 9 \), respectively. We suggest that they are lines between high Rydberg states of Ne v, with the same principal quantum numbers (level energies are given in NIST and Lapiere & Knystautas 1999). This is supported by our discovery of two high-\( n \) \( n = 8 \rightarrow 9 \) emission lines of Ne v at \( \lambda 5665 \)  \( \AA \) in some of our program stars (Fig. 5). Finally, the strong emission feature at \( \lambda 4945 \)  \( \AA \) was assigned to C v \( n = 7 \rightarrow 6 \). We think that it is the respective N v transition, as in the case of KPD 0005+5106. Nitrogen is definitely present in Sand 3, because the \( 3s \rightarrow 3p \) emission doublet at \( \lambda \lambda 4602, 4620 \)  \( \AA \) is present, as well as a prominent P Cygni profile of the N v resonance line. Thus, the entire optical emission line spectrum of Sand 3 and other [WCEs] can be explained without invoking ultrahigh-ionisation features.

Based on observations with the International Ultraviolet Explorer (IUE), several papers appeared in the literature claiming the presence of ultrahigh ionisation lines in the UV spectra of [WCE] stars. In Sand 3, as well as in NGC 5315 and NGC 6905, Feibelman (1996a, b) discovered the \( \lambda 1932 \)  \( \AA \) and 2977  \( \AA \) emission features that we have discussed above, and he also assigned them to O v. As in KPD 0005+5106 and RX J2117.1+3412 (Fig. 7), they probably are from Ne v.

Feibelman (1996b) even claimed the existence of an “O v sequence” of planetary-nebulae nuclei, mainly based on the supposed O v-iv identifications in PG1159 and [WCE] stars. Obviously, that idea must now be discarded. This is further corroborated by the fact that we detected the Ne v lines at \( \lambda 1162 \rightarrow 1165 \)  \( \AA \) in a [WCE] star too, namely in NGC 2371 (Fig. 6). The line strengths suggest \( T_{\text{eff}} \approx 150 000 \) K, which is, considering error limits, in good agreement with the result of a detailed analysis of FUSE and IUE spectra (135 000 K; Herald & Bianchi 2004). NGC 2371 was classified as [WC4] by Acker & Neiner (2003). These authors have based a new spectral classifi-
Fig. 5. First identification of two high-l \( n = 8 \to 9 \) emission lines of Ne \( \text{viii} \) (!). A strong and broad emission feature is seen in [WCE] stars at this location. We propose that it is due to these Ne \( \text{viii} \) transitions and not due to ultrahigh ionized O \( \text{vii} \). The labels denote the theoretical line positions using energy levels that are not precisely known.

Fig. 6. Identification of Ne \( \text{vii} \) lines in the [WCE] central star NGC 2371 (top) and comparison to the PG1159-type central star Longmore 4 (bottom).

6. Summary and conclusions

We have identified Ne \( \text{vii} \) absorption lines in FUSE spectra of PG1159 stars and a DO white dwarf. Line profile fits confirm that the PG1159 stars are the very hottest members of their spectral class \( (T_{\text{eff}} \geq 150,000 \text{ K}) \). We have shown that two well-known emission lines in the optical spectra of these stars are also from Ne \( \text{vii} \) and not, as previously thought, from O \( \text{vii} \). While Ne \( \text{vii} \) is thermally excited in the hot photospheres, the existence of O \( \text{vii} \) would require temperatures of the order of one million K, thus, an unknown process (e.g., shock zones in the wind) was invoked to explain these emission lines. This is no longer necessary.

We argue that these, and probably all, emission lines in early-type [WC] central stars that are usually assigned to ultrahigh-ionisation stages (O \( \text{vii} \)-\( \text{viii} \), C \( \text{v} \)) originate from Ne \( \text{vii} \)-\( \text{viii} \) and other thermally excited species.

The discovery of Ne \( \text{vii} \) lines in the central star Longmore 4 suggests that its \( T_{\text{eff}} \) is distinctively higher than previously thought (170,000 K instead of 120,000 K). To a smaller extent, the same tendency is seen for the PG1159 star K1-16 and the [WCE] star NGC 2371. This calls for a re-analysis of the complete spectra of these stars to see how the fit to other line features can be reconciled. The computation of model grids with varying parameters \( T_{\text{eff}} \), log \( g \), and abundances of the most important elements (He, C, O, Ne) will be necessary. Since the stars display prominent P Cygni line profiles, it could be advantageous to use expanding atmosphere models with the available spectral data. As already mentioned, the neon abundance cannot be determined from the newly discovered Ne \( \text{vii} \) lines because of the uncertainties with the line broadening theory. Instead, the strong P Cygni profile of the Ne \( \text{vii} \) at 973 Å line exhibited by these stars can be used as a sensitive tool (Bianchi & Herald 2007). Again, detailed parameter studies must be performed to find which neon abundance and which mass-loss rate yield good fits to this Ne \( \text{vii} \) line profile, while at the same time determining the \( T_{\text{eff}} \) of the model that is high enough to produce detectable Ne \( \text{vii} \) lines.

A surprising result of our investigation is the identification of Ne \( \text{vii} \) lines in the hottest known DO white dwarf, KPD 0005+5106. We conclude that its temperature is close to \( T_{\text{eff}} = 200,000 \text{ K} \) and, hence, that it is significantly hotter than hitherto thought \( (120,000 \text{ K}) \). As in the case of the PG1159 stars, the optical Ne \( \text{vii} \) emission lines were previously assigned to ultrahigh-ionized O \( \text{vii} \). The announcement of the discovery of a relatively soft X-ray emitting corona about KPD 0005+5106 seemed to support this assignment (Fleming et al. 1993). However, the analysis of a Chandra spectrum exhibiting flux in the range 20–80 Å proved that the soft X-rays stem from the photosphere of KPD 0005+5106 (Drake & Werner 2005). The deposition of the corona is now in accordance with the deposition of the ultrahigh-ionisation lines. A comprehensive re-analysis of all available data (Chandra, FUSE, HST, Keck) is required to tightly constrain the atmospheric parameters and metal abundances.

It remains to be seen how the extremely high effective temperature derived in our work relates to the observed hard X-ray emission at 1 keV \( (12 \text{ Å}) \) from KPD 0005+5106 (O’Dwyer et al. 2003). It might be that this emission is also of photospheric origin. After all, the entire spectral properties of KPD 0005+5106 could be deciphered as thermal photospheric radiation. In any case, the star has turned out to be by far the hottest known DO white dwarf.

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