Experimental Characterization of Unsteady Forces Triggered by Cavitation on a Centrifugal Pump

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Abstract
The paper presents an experimental campaign aimed at the characterization of the relationship between cavitation-induced instabilities and forces acting on the shaft relevant to space application turbopumps. The experiments have been carried-out on a 6-bladed, unshrouded centrifugal turbopump. The described apparatus allows for contemporary measurements of the flow instabilities and the force exchanged between the impeller and the shaft. Pressure fluctuations are frequency analyzed allowing for understanding the instability nature (axial, rotating) and their main characteristics (e.g. amplitude, rotating direction). The frequency content of the force components highlights a strong relationship of the z-component (along the rotating axis) with axial instabilities. On the other hand, rotating cavitation may involve force oscillations along all the three components leading to unwanted and dangerous fluctuating unbalances perpendicular to the rotating axis.

Keywords
Cavitation — Rotordynamic — Centrifugal Pump

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INTRODUCTION
Propellant feed turbopumps are a crucial component of all liquid propellant rocket engines due to the severe limitations associated with the design of dynamically stable, high power density machines capable of meeting the extremely demanding suction, pumping and reliability requirements of modern STSs (Space Transportation Systems).

The attainment of such high power/weight ratios is invariably obtained by running the impeller at the maximum allowable speed and low shaft torque. Current configurations in space applications are typically characterized by the presence of lighter, but also more flexible, shafts. Therefore, since the operation under cavitating conditions is tolerated, the turbopump is exposed to the onset of dangerous fluid dynamic and rotordynamic instabilities which may be triggered by cavitation phenomena. When designing a turbomachine, particularly if it has to operate at high rotational speeds, it is important to be able to predict the fluid-induced forces acting on the various components of the machine. The study of radial and rotordynamic forces on turbomachines components, by means of analytical/numerical and experimental approaches has been extensively carried out in the last 50 years by many researchers all around the world, as in the following refs. [1-7]. However, the experimental characterization of the influence of cavitation on these phenomena is still very poor even if it is a common operational condition in space application. Moreover, it is extremely important since the occurrence of cavitation drastically modifies the inertia of the fluid surrounding the impellers and, in turn, the critical speeds of the machine.

Fluid instabilities that develop in space inducers have been widely studied in the past, such in the following sources: [9-17]. However, a lack in centrifugal pumps experimental activities does not allow a complete understanding of the cavitation induced flow phenomena. In fact, few detailed experimental results can be found in literature, analyzing the instabilities and their interaction with the system even if rotating cavitation [18] and auto-oscillation [19] have been detected and studied in previous works.

The experimental characterization of the unsteady fluid forces/moments acting on space turbopump impellers as a consequence of the onset of the most dangerous types of cavitation-induced instabilities is here described and analyzed.

This is done by employing a new methodology, consisting in measuring simultaneously the forces acting on the shaft by means of a rotating dynamometer previously employed in past experimental activities [20-22] and the pressure field around the centrifugal pump at different locations by means of piezo-electric pressure transducers. This technique allows for comparing and evaluating the influence of pressure fluctuations on the unsteady force component that the pump exerts on the shaft, making possible to relate the nature of the unsteady flow phenomena interacting with the system by means of mechanical stresses.
1 ANALYSIS AND EXPERIMENTAL SETUP

1.1 Experimental apparatus

Figure 1 shows the general layout of the Cavitating Pump Rotordynamic Test Facility (CPRTF, see [23] for further information) at SITAEL S.p.A. (formerly ALTA S.p.A.). The facility is specifically intended for the experimental analysis of the relevant phenomena on space application turbopumps with particular attention on cavitation-related phenomena.

In order to relate the fluid instabilities with the unsteady forces acting on the pump, the experiments have been carried out by jointly using a rotating dynamometer (extensively used in past experimental activities, [24]), and a suitable number of piezoelectric pressure transducers (PCB Piezotronics, mod. S112A22) flush-mounted on the machine casing. The dynamometer has been placed downstream of the supports, between the pump and the rotor shaft, in order to avoid the influence of seals and bearing forces. The positions of the pressure transducers have been chosen in order to evaluate and identify the nature of the fluid instability (rotational/axial), thus different axial and azimuthal positions have been exploited. For this purpose, three different axial positions have been made available for the transducers on the pump casing (Figure 2 and Figure 4): upstream, midstream, and downstream (at the diffuser). With reference to Figure 4, the relative positions of the pressure taps have been designed in such a way that, at the same instant, number-like taps should face the through-flow of the same blade channel. The taps exploited during the present campaign are:

**Table 1.** Pressure taps exploited for the present campaign.

| Position     | Tap   |
|--------------|-------|
| Upstream     | 1, 2, 3, 4 |
| Midstream    | 1, 7  |
| Downstream   | 1, 2, 3, 7 |

**Figure 1.** The CPRTF lay-out in SITAEL S.p.A.

With reference to Figure 1 and Figure 2, the test setup includes:

- two electromagnetic flowmeters (Fisher-Rosemount, model 8732E, range: 0-100 l/s, accuracy 0.5% FS), for the measurement of the inlet and outlet flow rates;
- an absolute pressure transducer (\(p_{\text{up}}\)@\(\text{OPT1}\)) placed 6 diameters upstream w.r.t. the blade leading edge (GE, model Druck PMP1400, range 0÷1 bar, accuracy 0.25% FS) for the assessment of the inlet cavitation number (\(\sigma\));
- an absolute pressure transducer (\(p_2\)@\(\text{OPT2}\)) placed in the test chamber/section in order to measure the influence of the pressure on the axial force measured by the dynamometer. Indeed, there is a measured force related only to the pressure difference between the water pressure and the sealed dynamometer inside which is at atmospheric condition (GE, model UNIK5000, range 0÷6 bar, accuracy 0.1% FS);
- one differential pressure transducer (\(\Delta p\) @\(\text{OPT1 to OPT2}\)) which measures the pump pressure rise between the inlet station placed at 6 diameters upstream the blade leading edge and outlet station placed about two diameters downstream the blade trailing edge (GE, model UNIK5000, range 0÷5 bar, accuracy 0.1% FS);
- one temperature sensor PT100 (range 0÷100 °C, accuracy 0.5°C).

The test item employed in the present experimental campaign has been designed at SITAEL and it is a six-bladed radial pump (Figure 3) whose main characteristics are summarized in Table 2. The design is based on the reduced order model developed at ALTA S.p.A. [25, 26] which has been already applied to a similar radial impeller [27]. The clearance between the impeller and the casing is equal to 1 mm.

**Table 2.** Most relevant geometrical and operational parameters of the VAMPUFF pump.

| Parameter                              | Value            |
|----------------------------------------|------------------|
| Design flow coefficient @ \(r_2\)      | \(-\)            |
| \(\Phi_D\)                             | 0.120            |
| Number of blades                       | \([-\]\)         |
| \(N\)                                  | 6                |
| Outlet radius                          | \([mm]\)         |
| \(r_2\)                                | 105.00           |
| Inlet tip radius                       | \([mm]\)         |
| \(r_{T1}\)                             | 55.50            |
| Inlet hub radius                       | \([mm]\)         |
| \(r_{H1}\)                             | 31.42            |
| Axial length                           | \([mm]\)         |
| \(z_{H2}\)                             | 67.00            |
| Inlet tip blade angle                  | \([deg]\)        |
| \(\gamma_{T1}\)                       | 47.70            |
| Inlet backswep angle                   | \([deg]\)        |
| \(\chi_1\)                            | 0                |
| Diffuser outlet radius                 | \([mm]\)         |
| \(r_3\)                               | 126              |
| Design Rotating Speed                  | \([rad/s]\)      |
| \(\Omega\)                             | 1500            |
| \([\text{rpm}]\)                      | (1500)          |
| Design volumetric flow rate            | \([m^3/s]\)      |
| \(Q_{\text{des}}\)                    | 0.022            |
| Mean blade height                      | \([mm]\)         |
| 18.84                                  |
| Tip solidity                           | \([-\]\)         |
| \(\sigma_T\)                          | 2.64             |
| Incidence tip angle @ design           | \([deg]\)        |
| \(\alpha\)                            | 18.90            |
| Outlet tip blade mean angle            | \([deg]\)        |
| \(\gamma_{T2}\)                       | 89.94            |
| Outlet tip backswep angle              | \([deg]\)        |
| \(\chi_2\)                            | 64.00            |
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1.1 Pumping Performance

The tests are performed at room temperature \( T \approx 20^\circ C \), while the inlet pressure is set well above the vapor pressure in order to avoid cavitation inception within the pump. Referring to the pump geometric in Table 2, the parameters used to characterize the performance are:

- Total Head Coefficient \( \Psi_T = \frac{\Delta p}{\rho \Omega^2 r^2} \)
- Flow Coefficient \( \Phi = \frac{Q}{\Omega r^2} \)
- Reynold Number \( Re = \frac{2\Omega r^2}{v} \)

1.2 Cavitating Performance

Steady-state cavitating experiments are performed by maintaining a fixed inlet pressure condition during the experiment. Continuous cavitating experiments are performed by slowly lowering the inlet pressure in order to reach different cavitating regimes. The acquired signals of each continuous experiment are divided into sub-blocks where the operating conditions and relevant phenomena do not change significantly. The main parameters used to characterize the cavitating performance are:

- Head Coefficient \( \Psi = \frac{\Delta p}{\rho \Omega^2 r^2} \)
- Cavitation Number \( \sigma = \frac{(p_{in} - p_v)}{0.5 \rho \Omega^2 r^2} \)

1.3 Flow Instabilities and Unsteady Fluid Forces

The pressure fluctuations of the flow through the pump triggered or not by cavitation have been investigated during the cavitating experiments. Cervone, et al., and d'Agostino, et al. [16, 28] report the procedures already exploited by the authors and here just summarized. As for the cavitating performance evaluation, the signals acquired by means of the piezoelectric transducers located on the pump casing and diffuser are divided into sub-blocks. Per each sub-block, the signals are fast-Fourier transformed in order to obtain the frequency content of oscillating phenomena. At each frequency of interest, the analysis of the cross-spectrum phase of signals from transducers placed at the same axial location but at different azimuthal positions allows for understanding the physical nature of the phenomenon. In fact, oscillating axial phenomena involve signals characterized by zero-phase cross-spectrum. On the other hand, rotating phenomena lead to cross-spectrum phases \( \varphi \) proportional to the circumferential spacing \( \Delta \theta \) between transducers and proportional to the phenomenon number of lobes \( n \). For instance, for a generic cross-spectrum from the transducer “1” to the transducer “2”, it would be:
where negative values of \( \varphi_{12} \) imply a phenomenon moving from transducer “1” to transducer “2”. Aliasing in the azimuthal direction is eliminated by comparison of the cross-correlations of pressure signals from transducers with different angular separations. The analysis is validated based on the value of the coherence function between signals obtained from two generic transducers \((x, y)\) \( \gamma_{xy} = \left| \frac{S_{xy}(f)}{S_{xx}(f)S_{yy}(f)} \right| \), where \( S \) is either the cross-spectrum or the auto-spectrum of the referred transducers \((x, y)\) in the subscripts (see [29] for further information). Here, only phenomena with values of \( \gamma_{xy} \) greater than 0.95 have been considered.

In order to understand the frequency content of the signals recorded through the dynamometer, it is useful to consider and compare the effect of different types of unsteady forces on the dynamometer rotating frame \((xyz)\) and on the absolute fixed one \((XYZ)\)-where \(x\) and \(Z\) are directed along the rotating axis. Therefore, the force evaluated in the two reference frames has the same \(z\)-component as well as the same total force acting on the plane \(X - Y\) \((x - y)\).

Consideration of Figure 5, the schematic proposed involves a fluctuating phenomenon generating a force with constant direction \((\theta_0)\) w.r.t. the fixed reference frame, a mean value \(F_0\) and an oscillating amplitude: \(F_1 = F_0 + \bar{F}_1 \cos(\omega_1 t)\). Furthermore, the schematic considers a constant amplitude force \(F_2\) rotating w.r.t. the fixed reference frame with an angular velocity \(\omega_1\). The force components in the rotating frame and in the fixed one are:

\[
F_x = F_0 \cos(\theta_0) + \bar{F}_1 \cos(\omega_2 t) \cos(\theta_0) + \bar{F}_2 \cos(\omega_1 t)
\]

\[
F_y = F_0 \sin(\theta_0) + \bar{F}_1 \cos(\omega_2 t) \sin(\theta_0) + \bar{F}_2 \sin(\omega_1 t)
\]

\[
F_z = \bar{F}_1 \cos[(\omega_2 - \Omega) t + \theta_0] + \frac{\bar{F}_1}{2} \cos[(\omega_2 + \Omega) t - \theta_0] + \bar{F}_2 \cos(\Omega t - \theta_0)
\]

\[
F_y = \bar{F}_1 \sin[(\omega_2 - \Omega) t + \theta_0] - \frac{\bar{F}_1}{2} \sin[(\omega_2 + \Omega) t - \theta_0] + \bar{F}_2 \sin(\Omega t - \theta_0)
\]

Therefore, a frequency analysis of the force components would show 2 frequencies \((\omega_1, \omega_2)\) in the fixed reference frame. On the other hand, the same analysis in the rotating frame would show multiple frequencies shifted from the above-mentioned ones by \(\pm \Omega\).

### Table 3. Relevant parameters for the cavitating experiments.

| Sampling Frequency \((f_s)\) | 5000 Hz |
| Temperature | 20 °C |
| Experiment Duration (cont.) | 240 s |

### Cavitation Performance, \(T = 20^\circ \text{C}\)

![Figure 7. VAMPUFF cavitating performance at different flowrates \((T = 20^\circ\text{C}; \Omega = 1750 \text{ rpm})\).](imageURL)

### 2 RESULTS AND DISCUSSION

#### 2.1 Pumping Performance

Figure 6 shows the pumping performance obtained at three different rotating speeds. The data overlapping confirms the Reynolds independence of the tests (fully turbulent flow).

![Figure 6. VAMPUFF pumping performance, \(T = 20^\circ \text{C}\).](imageURL)

#### 2.2 Cavitating Performance

Figure 7 reports the VAMPUFF cavitating performance obtained at different flow coefficients \(\Phi\). Per each flow coefficient, solid markers with black borders show results obtained from steady-state experiments, while blank markers report continuous experiments results. The overlapping of the data from the two different procedures confirms the goodness of the continuous data which can reliably represent the cavitating performance. The rotational speed \(\Omega = 1750 \text{ rpm}\) allows for measurable cavitational and forces without exceeding the maximum capability of the dynamometer.
2.3 Flow Instabilities and Unsteady Fluid Forces

The systematic evaluation of the flow instabilities during the cavitating experiments highlighted the presence of rotating and axial phenomena for different flowrates, as briefly summarized in Table 4. In the table, at the same flowrate, oscillating phenomena showing analogous characteristics (rotating with 1 lobe, axial, etc.) are grouped together regardless the operating regimes ($\sigma$, frequency).

| $\Phi$ ($\Phi/\Phi_D$) | Identified Phenomena          |
|-----------------------|-------------------------------|
| 0.096 (0.8)           | Rotating (1 lobe and 2 lobes); axial |
| 0.108 (0.9)           | Rotating (1 lobe and 2 lobes); axial |
| 0.120 (1.0)           | Rotating (1 lobe); axial       |
| 0.132 (1.1)           | Rotating (1 lobe); axial       |
| 0.144 (1.2)           | Axial                         |

The presence of cavitation instabilities leads to unwanted forces on the shaft even at $\Phi = \Phi_D$. However, at design condition the intensity of the detected oscillating forces is minor than at lower flowrates. For this reason, the present study only focuses on $\Phi = 0.108$ ($\Phi/\Phi_D = 0.9$) in order to show some of the typical pressure-force spectra relationships found at all the tested regimes.

Figure 8 reports the cavitating performance at a nominal $\Phi_N = 0.108$ as well as the flow coefficient behavior during the inlet pressure decay. At the end of the experiment the massive presence of cavitation leads to the breakdown which, in turn, leads to the flow coefficient drop.

Figure 8. Cavitating performance and the flow coefficient evolution, nominal $\Phi_N = 0.108$ (90% $\Phi_D$).

Figure 9. Frequency energy content $[Pa^2 \cdot s]$ of pressure transducers placed upstream (left), midstream (center), and downstream (right).
The flow instabilities illustrated in this paper are basically connected with the presence of cavitation. According to [30], cavitation inception usually starts in the tip vortex generated at the blade inlet. Therefore, it is of no surprise that the major part of the found phenomena is clearly visible by the upstream pressure transducers, becoming less visible while moving downstream. On the other hand, axial phenomena effectively propagate from the upstream to the downstream and vice versa becoming clearly visible in both stations.

Figure 9 reports the energy frequency content of three pressure transducers as representative of the three different stream locations: upstream, midstream, and downstream (according to Figure 2 and Figure 4) versus the cavitation number \( \sigma \). Per each frequency, the frequency energy content \( E_f \) is directly related to the amplitude \( A_f \) of the acting oscillating phenomenon as follows:

\[
E_f = \frac{A_f^2 N_s}{4 f_s}
\]

where \( N_s \) is the number of samples considered for the FFT while \( f_s \) is the sampling frequency exploited during the experiment.

While the blade passage frequency is clearly visible at \( 6\Omega \) per each station, the other relevant phenomena are generally of major interest at a single station. In order to understand the physical nature of such oscillating phenomena, Figure 10, Figure 11, and Figure 12 report the phase of the cross-spectrum of the upstream, midstream, and downstream pressure transducers, respectively. The figures show only the phenomena of major interest characterized by an amplitude \( A_f \geq 100 \text{ Pa} \) and a coherence \( \gamma_{xy} \geq 0.95 \), focusing on the range of \( f/\Omega = 0 + 4 \) where interesting phenomena have been found. The phase values allow for understanding the nature of the phenomenon [28, 31]. The main outcomes of this analysis are summarized in Table 5.

**Table 5. Summary of the found instabilities at \( \Phi_N = 0.108 \).**

| ID | Freq. Range \( f/\Omega \) \( (f [Hz]) \) | \( \sigma \) range | Characteristics | Major Station |
|----|----------------------------------|----------------|----------------|---------------|
| A1 | 0.38 - 0.1 \( (11 - 3) \)        | 0.16 - 0.12    | Axial          | Up            |
|    | 2.74 - 2.57 \( (80 - 75) \)      | 0.2 - 0.16     | Rotating 1 lobe| Up            |
|    | 2.85 - 2.33 \( (83 - 68) \)      | 0.16 - 0.11    | Rotating 2 lobes| Mid           |
| A2 | 3.22 - 3.81 \( (94 - 111) \)     | 0.16 - 0.10    | Axial          | Down          |
|    | 3.26 - 3.46 \( (95 - 101) \)     | 0.19 - 0.14    | Rotating 1 lobe| Mid           |

**Figure 10.** Cross-spectrum phase of the upstream pressure transducers, \( \Phi = 0.108 \) (90%\( \Phi_D \)).

**Figure 11.** Cross-spectrum phase of the midstream pressure transducers, \( \Phi = 0.108 \) (90%\( \Phi_D \)).
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Figure 12. Cross-spectrum phase of the downstream pressure transducers, $\Phi = 0.108$ (90% $\Phi_D$).

Figure 13 shows the frequency energy content of the force components measured by the dynamometer with a force amplitude oscillation greater than 0.5 N. In particular, the figure reports the force components independent of the chosen reference frame (i.e. rotating or fixed), as described previously.

RC2 and RC3 do not generate relevant effects on the forces sensed by the dynamometer, therefore they won’t be further discussed. On the contrary, some of the instabilities reported in Table 5 clearly lead to unwanted oscillations of the force acting on the shaft. In particular, the two axial instabilities A1 and A2 generate oscillations of $F_z$ for corresponding values of $\sigma - f$, while they are not visible at all on a plane perpendicular to the rotational axis ($XY, xy$).

Like A1 and A2, the pressure distribution connected with the rotating cavitation-induced instability RC1 generates a fluctuating component on the rotational axis. However, RC1 also leads to a fluctuating component on the plane $XY$ ($xy$) at the same operating regimes ($\sigma$) and frequencies, with an intensity directly connected to the corresponding energy value.

In order to understand the RC1 effects on the plane $XY$ ($xy$), it is useful to analyze the frequency content of the force component $F_X$ and $F_Z$ defined in the fixed frame and in the rotating one, respectively (Figure 14). Let’s consider the schematic proposed in Figure 5 and in the following. The frequency energy content of $F_{xy}$ is influenced only by the amplitude and the acting frequency of $\tilde{F}_1$. The pressure distribution due to the presence of cavitation in the form of RC1 generates a rotating force imbalance with a rotational velocity corresponding to the rotational velocity of the phenomenon itself, which is given by the pressure transducer analysis. Moreover, when this pressure distribution interacts with the (static) volute tongue at the impeller exit, it may lead to an oscillating unbalanced force at the same frequency as the phenomenon itself. The above considerations may be summarized as $\omega_2 = \omega_1 = 2\pi f_{RC1}$, therefore:

$$F_X = F_0 \cos(\theta_0) + [F_1 \cos(\theta_0) + F_2] \cos(\omega_1 t)$$

which is coherent with the frequency content in Figure 14. Moreover, the corresponding force on the rotating frame is given by:

$$F_x = \frac{\tilde{F}_1}{2} \cos[(\omega_1 - \Omega) t + \theta_0] + F_2 \cos[(\omega_1 - \Omega) t]$$

$$+ \frac{F_1}{2} \cos[(\omega_1 + \Omega) t - \theta_0] + F_0 \cos(\Omega t - \theta_0)$$

where there are only three acting frequencies (as shown in Figure 14, center):

1. $(\omega_1 - \Omega)/2\pi$, whose intensity depends on the combination of $F_1, F_2, \theta_0$;
2. $(\omega_1 + \Omega)/2\pi$;
3. $\Omega$.
Figure 13. Frequency energy content \( [N^2 \cdot s] \) of the force along the axis \( (F_z) \), its perpendicular component \( (F_{xy}) \), and the total force \( (F_{xyz}) \) acting on the dynamometer.

Figure 14. Frequency energy content \( [N^2 \cdot s] \) of the force component along the X-axis in the absolute fixed reference frame \( (F_x, \text{left}) \), along the x-axis in the rotating reference frame \( (F_x) \) and its perpendicular component \( (F_{xy}) \) acting on the dynamometer.
3 CONCLUSIONS

The experimental characterization of pressure fluctuations on a centrifugal pump for various operating regimes has been here presented and discussed. The paper focuses on the unsteady forces acting on the driving shaft which are triggered by such fluctuations. Although the paper presents data coming from a specific regime in terms of flow coefficient ($\Phi = 0.9 \Phi_p$), similar behaviors have been detected for analogous phenomena observed during other regimes. In particular, the presence of rotating phenomena may generate fluctuations of the force in the plane perpendicular to the rotational axis as well as on the rotational axis itself. On the perpendicular plane, the frequency content associated with the generated force suggests that the force can be approximated by the sum of a component with a fixed direction and fluctuating intensity together with a purely rotating component. The frequency of both components is the same as for the source phenomenon. A key role in such behavior is most likely played by the single-tongue volute whose intrinsic asymmetry interacts with the rotating phenomena leading to force amplitude fluctuations. At $\Phi = \Phi_p$ the outlet flow has a nominal zero incidence angle with the tongue, which is most likely the reason why the force triggered by cavitation instabilities is also reduced. Thus, confirming the goodness of the pump design approach described in [25]. However, experimental campaigns are needed to better understand the interaction between the presence of cavitation instabilities and forces in presence of different shaped volutes and diffusers. Furthermore, the paper highlights that special attention should be paid when the pump design includes any stator (e.g. a vaned diffuser) especially when the operating conditions may include cavitation.

On the other hand, axial phenomena lead to fluctuations of the force component directed along the rotational axis only.

To date, it is not clear which eventual unstable forces would arise with an inducer ahead of a centrifugal stage (typical in space application) under the presence of cavitation instabilities. However, the present results suggest that the massive presence of such instabilities for an inducer (if compared to a centrifugal stage) may lead to possible dangerous conditions when the inducer is coupled with static elements typical of centrifugal stages (e.g. the volute tongue). Further investigations are needed in order to better clarify these aspects.

4 NOMENCLATURE

| Symbol | Description                           |
|--------|---------------------------------------|
| latin  |                                       |
| A      | Axial                                 |
| $A_f$  | Fluctuation amplitude                 |
| $E_f$  | Fluctuation energy                    |
| $F$    | Force intensity fluctuation           |
| $F$    | Force                                 |
| $N$    | Blade Number                          |
| $Q$    | Volumetric Flow Rate                  |
| $RC$   | Rotating Cavitation                   |
| $Re$   | Reynolds                              |
| $S$    | Power Spectrum Density                |
| $T$    | Temperature                           |
| $f$    | Frequency                             |
| $n$    | Flow instability lobes                |
| $p$    | Pressure                              |
| $r$    | Radius                                |
| $z$    | Axial Length, rotating axis           |

Greek symbols:
- $\Delta$: Variation
- $\Phi$: Flow Coefficient
- $\Psi$: Head Coefficient
- $\Omega$: Rotational Speed
- $\gamma_{xy}$: Coherence between sensors $x$ and $y$
- $\nu$: Fluid Kinematic Viscosity
- $\rho$: Fluid Density
- $\sigma$: cavitation number
- $\omega$: Force angular frequency
- $\varphi$: Cross-spectrum phase
- $\theta$: Force initial phase
- $\vartheta$: Azimuthal direction

Subscripts:
- $D$: Design
- $N$: Nominal
- $T$: Total
- $c$: chamber
- $in$: Inlet
- $v$: vapor
- $xyz$: Rotating frame
- $XYZ$: Fixed frame

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