Improving energy harvesting in excited Duffing harvester device using a delayed piezoelectric coupling

Z. Ghouli, M. Hamdi, M. Belhaq

1 Faculty of Sciences Ain Chock, University Hassan II-Casablanca, Morocco
2 FST-Al Hoceima, Mohamed First University Oujda, Al Hoceima, Morocco

Abstract. The present work examines the influence of time delay introduced in the piezoelectric circuit of an excited Duffing harvester device with hardening stiffness on the vibration and voltage amplitudes. Specifically, we seek to exploit a delayed electrical circuit of the harvester to enhance its performance. We consider the case of a monostable system and we use a perturbation technique to approximate the periodic response and the corresponding voltage amplitude near the principal resonance. It is shown that for appropriate values of delay amplitude, the energy harvesting performance is improved over a certain range of coupling parameters and excitation frequencies. Numerical simulation is conducted to support the analytical predictions.

1 Introduction

A monostable piezoelectric nonlinear energy harvester (EH) device consisting of a Duffing-type harvester with hardening stiffness coupled to a delayed piezoelectric circuit is considered in this paper. The objective is to examine the influence of time delay in the piezoelectric circuit on the amplitude of the response and the voltage amplitude.

Recently, a delayed Duffing-type monostable harvester subject to a harmonic excitation and coupled to a piezoelectric circuit was studied in the case where the time delay is introduced only in the mechanical subsystem [1]. The optimal performance of the harvester device in terms of time delay parameters was investigated and it was shown that the induced large-amplitude quasiperiodic vibrations can be used to extract energy over broadband of excitation frequencies away from the resonance.

Here we examine the influence of time delay on EH in the case where the time delay is introduced in the electrical circuit. To show the contribution of introducing time delay in the electrical circuit in improving EH performance, a comparison is carried out with the case where time delay is absent [2].

In the next section we present the harvester system and we derive approximation of the periodic response and the voltage amplitude using the multiple scales method [3]. In Section 3 the influence of different system parameters of the harvester device on the EH performance is analyzed in the presence and absence of time delay. A summary of the results is given in the concluding section.

2 Harvester system, vibration response and voltage amplitude

Consider an excited Duffing oscillator coupled to an electrical circuit through a delayed piezoelectric device as shown in the schematic presented in Fig. 1.

The governing equations for this harvester can be written in the dimensionless form as

\[\ddot{x}(t) + 2\mu \omega_n \dot{x}(t) + \omega_n^2 x(t) + \alpha x(t)^3 - \chi v(t) = f \cos(\omega t) \quad (1)\]

\[\dot{v}(t) + \lambda v(t) + \kappa \dot{x}(t) = \beta \dot{v}(t - \tau) \quad (2)\]

where \(x(t)\) is the relative displacement of the rigid mass \(m\), \(v(t)\) is the voltage across the load resistance, \(\omega_n\) is the natural frequency of the system, \(\mu\) is the mechanical damping ratio, \(\alpha\) is the stiffness parameter, \(\chi\) is the piezoelectric coupling term in the mechanical attachment, \(\nu\) is the piezoelectric coupling term in the electrical circuit, \(\lambda\) is the reciprocal of the time constant of the electrical circuit, \(\beta\) and \(\tau\) are, respectively, the feedback gain and time delay in the electrical circuit, and \(f, \omega\) are, respectively, the amplitude and the frequency of the excitation. To investigate the influence of time delay in the electric circuit on the performance of the harvester, we approximate the response of the system near the primary resonance by introducing the resonance condition \(\omega = \omega_n + \sigma\) where \(\sigma\) is a detuning parameter.

The method of multiple scales [3] is implemented by introducing a bookkeeping parameter \(\epsilon\) and scaling parameters as \(\mu = \epsilon \hat{\mu}, \alpha = \epsilon \hat{\alpha}, \chi = \epsilon \hat{\chi}, \sigma = \epsilon \hat{\sigma}, f = \epsilon \hat{f}\).
Thus, Eqs. (1) and (2) take the form
\[ \ddot{x}(t) + \omega^2_n x(t) = -2\mu \omega_n \dot{x}(t) - \ddot{\alpha} x(t)^3 + \ddot{\chi} v(t) + f \cos(\omega t) \] (3)
\[ \ddot{v}(t) + \lambda \dot{v}(t) + \kappa \dot{x}(t) = \beta_\omega (t - \tau) \] (4)
A solution to Eqs. (3) and (4) is given by
\[ x(t) = x_0 (T_0, T_1) + \epsilon x_1 (T_0, T_1) + O(\epsilon^2) \] (5)
\[ v(t) = v_0 (T_0, T_1) + \epsilon v_1 (T_0, T_1) + O(\epsilon^2) \] (6)
where \( T_0 = t \) and \( T_1 = t + \epsilon t \). In terms of the variables \( T_1 \), the time derivatives become \( \frac{d}{dt} = D_0 + \epsilon D_1 + O(\epsilon^2) \) and \( \frac{d^2}{dt^2} = D_0^2 + \epsilon^2 D_1^2 + 2 \epsilon D_0 D_1 + O(\epsilon^2) \) where \( D_1^j = \frac{\partial}{\partial T_j} \). Substituting (5) and (6) into (3) and (4) and equating coefficients of like powers of \( \epsilon \), we obtain up to the second order the systems
\[ D_0^2 x_0 + \omega^2_n x_0 = 0 \] (7)
\[ D_0 x_1 + \lambda x_1 + \kappa D_0 x_0 = \beta_\omega v_0 \] (8)
\[ D_0^2 x_1 + \omega^2_n x_1 = -2 \mu D_0 x_0 - 2 \mu \omega_n D_0 x_0 - \ddot{\alpha} x_1^3 + \ddot{\chi} v_0 + f \cos(\omega_0 T_0 + \theta T_1) \] (9)
\[ D_0 v_1 + \lambda v_1 = -D_1 v_0 - \kappa D_0 x_1 - D_1 x_0 + \beta_\omega v_1 \] (10)
Up to the first order the solution reads
\[ x_0(T_0, T_1) = A(T_1) e^{i\omega_0 T_0} + \bar{A}(T_1) e^{-i\omega_0 T_0} \] (11)
\[ v_0(T_0, T_1) = \frac{-ki_0 \omega_n A(T_1)}{\lambda - i\omega_n - \beta_\omega e^{-i\theta} e^{-i\omega_0 T_0}} + \frac{ki_0 \omega_n \bar{A}(T_1)}{\lambda - i\omega_n - \beta_\omega e^{i\theta} e^{i\omega_0 T_0}} \] (12)
where \( A(T_1) \) and \( \bar{A}(T_1) \) are unknown complex conjugate functions. Substituting Eqs. (11) and (12) into (9) and (10) and eliminating the secular terms, one obtains
\[ -2i\omega_n (D_1 A) - 2i\mu \omega_n A - 3i\alpha A^2 - \frac{ki_0 \omega_n A}{\lambda - i\omega_n - \beta_\omega e^{-i\theta}} + \frac{fi_0}{2} e^{i\theta T_1} = 0 \] (13)
Expressing \( A = \frac{1}{2} \omega_\theta e^{i\theta} \) where \( a \) and \( \theta \) are the amplitude and the phase of the modulation, we obtain up to the first order the modulation equations
\[ \begin{cases} \frac{da}{dt} = S_1 a + S_4 \sin(\gamma) \\ \frac{d\gamma}{dt} = S_2 a + S_3 a^3 + S_4 \cos(\gamma) \end{cases} \] (14)
where \( S_j (j = 1, ..., 4) \) are given in Appendix and \( \gamma = \delta T_1 - \theta \). The steady-state response of system (14), corresponding to periodic oscillations of Eqs. (3) and (4), are determined by setting \( \frac{da}{dt} = \frac{d\gamma}{dt} = 0 \). Eliminating the phase, we obtain the following sixth-order algebraic equation in \( a \)
\[ (S_1 a)^2 + (S_2 a + S_3 a^3)^2 = S_4^2 \] (15)
which can be written in the form
\[ B a^6 + C a^4 + D a^2 + E = 0 \] (16)
where \( B = S_2^2, C = 2S_1 S_3, D = S_1^2 + S_3^2 \) and \( E = -S_4^2 \). The discriminant of Eq. (16) reads
\[ \Delta = \frac{p^4}{27} + \frac{q^2}{4} \] (17)
where \( P = \frac{D}{a} \) and \( Q = \frac{C}{a} \). Equation (16) has three real positive roots if \( \Delta \) is negative and one positive root if \( \Delta \) is positive.

The solution given by Eqs. (11) and (12) can be written as \( x_0(T_0, T_1) = a \cos(\omega_0 T_1 + \theta) \) and \( v_0(T_0, T_1) = V \cos(\omega_0 T_1 + \theta + \arctan \left( \frac{\lambda - i\omega_n - \beta_\omega e^{-i\theta} e^{-i\omega_0 T_0}}{A(T_1)} \right)) \) such that the condition \( \omega_0 + \beta \sin(\omega_0 T_1) \neq 0 \) is satisfied. Moreover, the voltage amplitude \( V \) is given by
\[ V = \frac{\omega_n}{\sqrt{\left(\lambda - \beta_\omega \cos(\omega_0 T_1)\right)^2 + \left(\omega_n + \beta_\omega \sin(\omega_0 T_1)\right)^2}} \] (18)

### 3 Main results

Next, the influence of different parameters of the system on vibration and voltage amplitudes is examined using Eqs. (15) and (18), respectively.

Figures 2a and 2b show, respectively, the variation of the amplitude of the periodic vibrations and the voltage \( V \) versus \( \sigma \) for the excitation amplitude \( \beta = 0.1 \). The undelayed case (\( \beta = 0 \)) is presented by a grey line and the delayed one (\( \beta 
eq 0 \)) is indicated by a black line. The analytical prediction (solid lines) is compared to results obtained by numerical simulation using dde23 algorithm [4] for the delayed case and the method of Runge Kutta of order 4 for the undelayed case. It can be observed from the figures that the EH performance is improved by introducing the time delay in the electric circuit, especially near the resonance peak.

Figure 3 shows the variation of the amplitude of the periodic response and the voltage amplitude \( V \) versus \( \sigma \) for larger value of the forcing amplitude \( \beta = 1.5 \). The grey line corresponds to the undelayed case (\( \beta = 0 \)) and the black one is obtained in the delayed one. As in the previous case (Fig. 2 for \( \beta = 0.1 \)) the voltage amplitude is increased around the resonance peak in the case of delayed circuit.

Similarly, the variation of the amplitude of vibration and voltage versus the piezoelectric coupling term in the electrical circuit \( \kappa \) are depicted in Fig. 4 showing the range of parameter \( \kappa \) where a substantial EH performance is achieved (Fig. 4b). The variation of the amplitudes of vibration and voltage as a function of the piezoelectric coupling term in the mechanical interaction \( \chi \) are shown in Fig. 5 for \( \beta = 0 \) (undelayed circuit, grey line) and \( \beta = \lambda \) (delayed circuit, black line). Inspection of Fig. 5 shows that for values of \( \sigma \) taken relatively away from the resonance, the range of \( \chi \) where a better performance of EH is obtained (Fig. 5b).

### 4 Conclusion

We have studied the EH performance in a monostable Duffing forced oscillator with hardening stiffness coupled to a delayed piezoelectric harvester device. The analysis was carried out near the primary resonance and a perturbation method was performed to obtain approximation of vibration and voltage amplitudes. The influence of time delay introduced in the piezoelectric subsystem on the EH performance of the harvester was studied. Results showed that
in the presence of time delay in the electrical circuit the voltage amplitude increases at certain range of coupling parameters especially in the vicinity of the resonance peak. In other words, by tuning the delay amplitude in the electric circuit, it is possible to optimize the power output performance over certain range of the piezoelectric coupling terms.

5 Appendix

\[ S_1 = -\mu \omega_n - \frac{\chi \kappa (\lambda - \beta \cos(\omega_n \tau))}{2[\lambda - \beta \cos(\omega_n \tau)]^2 + (\omega_n + \beta \sin(\omega_n \tau))^2]} \]

\[ S_2 = \sigma = -\frac{\chi \kappa (\omega_n + \beta \sin(\omega_n \tau))}{2[\lambda - \beta \cos(\omega_n \tau)]^2 + (\omega_n + \beta \sin(\omega_n \tau))^2]} \]

References

1. Z. Ghouli, M. Hamdi, F. Lakrad, M. Belhaq, Quasiperiodic energy harvesting in a forced and delayed Duffing harvester device, J. Sound Vib. 407 (2017) 271-285.
2. A. Erturk, D.J. Inman, Piezoelectric Energy Harvesting, John Wiley, 2011.
3. Nayfeh, A.H., Mook, D.T.: Nonlinear Oscillations. Wiley, New York, 1979.
Fig. 4. Vibration (a) and voltage (b) amplitudes vs $\kappa$ for $\sigma = 2, \chi = 5, f = 0.5, \omega_n = 1, \mu = 0.5, \alpha = 5, \lambda = 0.05$, and $\tau = 6.2$. Black (grey) line for delayed (undelayed) electric circuit $\beta = \lambda$ ($\beta = 0$). Analytical prediction (solid lines) and numerical simulation (Black (grey) circles for delayed (undelayed) electric circuit).

4. L.F. Shampine, S. Thompson, Solving delay differential equations with dde23. PDF available on-line at http://www.radford.edu/~thompson/webddes/tutorial.pdf (2000).

Fig. 5. Vibration (a) and voltage (b) amplitudes vs $\chi$ for $\sigma = 2, \kappa = 0.5, f = 0.5, \omega_n = 1, \mu = 0.5, \alpha = 5, \lambda = 0.05$, and $\tau = 6.2$. Black (grey) line for delayed (undelayed) electric circuit $\beta = \lambda$ ($\beta = 0$). Analytical prediction (solid lines) and numerical simulation (Black (grey) circles for delayed (undelayed) electric circuit).