STAR CLUSTER PROPERTIES IN TWO LEGUS GALAXIES COMPUTED WITH STOCHASTIC STELLAR POPULATION SYNTHESIS MODELS

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ABSTRACT

We investigate a novel Bayesian analysis method, based on the Stochastically Lighting Up Galaxies (slug) code, to derive the masses, ages, and extinctions of star clusters from integrated light photometry. Unlike many analysis methods, slug correctly accounts for incomplete initial mass function (IMF) sampling, and returns full posterior probability distributions rather than simply probability maxima. We apply our technique to 621 visually confirmed clusters in two nearby galaxies, NGC 628 and NGC 7793, that are part of the Legacy Extragalactic UV Survey (LEGUS). LEGUS provides Hubble Space Telescope photometry in the NUV, U, B, V, and I bands. We analyze the sensitivity of the derived cluster properties to choices of prior probability distribution, evolutionary tracks, IMF, metallicity, treatment of nebular emission, and extinction curve. We find that slug’s results for individual clusters are insensitive to most of these choices, but that the posterior probability distributions we derive are often quite broad, and sometimes multi-peaked and quite sensitive to the choice of priors. In contrast, the properties of clusters are insensitive to most of these choices, but that the posterior probability distributions we derive are often quite broad, and sometimes multi-peaked and quite sensitive to the choice of priors. We also compare our results from slug to those derived with a conventional non-stochastic fitting code, Yggdrasil. We show that slug’s stochastic models are generally a better fit to the observations than the deterministic ones used by Yggdrasil. However, the overall properties of the cluster populations recovered by both codes are qualitatively similar.

Key words: galaxies: individual (NGC 628, NGC 7793) – galaxies: star clusters: general – methods: data analysis – methods: statistical – techniques: photometric

Supporting material: machine-readable table

1. INTRODUCTION

Star clusters represent a scale of star formation intermediate between individual stars and entire galaxies. They likely represent the gravitationally bound peaks of a continuous distribution of star formation across length and time scales. For this reason, the study of their properties, and any variation in those properties as a function of galactic environment, is critical to an overall theoretical understanding of the star formation process (see Portegies Zwart et al. 2010; Longmore et al. 2014; Krumholz 2014 for recent reviews).

While the study of star clusters is an old topic, except in the Milky Way and the very closest external galaxies, much of this work has by necessity focused on the relatively massive clusters, ∼10^5.5 M_⊙ or more (e.g., Zhang & Fall 1999; Larsen & Richtler 2000; Bik et al. 2003; de Grijs et al. 2003b; Fall et al. 2005; Goodwin & Bastian 2006; Larsen 2009). This has created significant challenges for theoretical interpretation, for two reasons. The first is simply statistics: as discussed in more detail below, observed star cluster mass functions appear to be well-fit by power laws of the form dN/dM ∝ M^(-2) (e.g., Williams & McKee 1997; Zhang et al. 1999; Larsen 2002; Bik et al. 2003; de Grijs et al. 2003b; Goddard et al. 2010; Bastian et al. 2011; Fall & Chandar 2012; Fouesneau et al. 2012). Such a mass function implies that clusters in the mass range 10^2.5−10^3.5 M_⊙ are nearly ten times as numerous as those in the range ∼10^3.5−10^5.5 M_⊙ that is typically studied, so a sample consisting only of massive clusters necessarily has much less statistical power than one probing to lower masses. The second is that theoretical models for the evolution and disruption of star clusters, either due to stellar feedback or due to environmental influence, make some of their strongest predictions for the effects of cluster mass at masses below this range (e.g., Lamers et al. 2005; Parmentier et al. 2008; Fall et al. 2010; Kruisjes 2012; Kruisjes et al. 2012).

In the last five years a number of groups have made a concerted effort to extend the star cluster data set to lower...
masses and a broader range of galactic environments, a necessary step if we are to differentiate between models. Doing so is challenging both observationally and theoretically. Observationally, ground-based studies probing to lower masses are mostly limited by resolution to the Milky Way (Borissova et al. 2011) and the Magellanic Clouds (Hunter et al. 2003; Rafelski & Zaritsky 2005). Extending this sample is one of the primary goals of two ongoing Hubble Space Telescope (HST) programs: the Panchromatic Hubble Andromeda Treasury (PHAT; Dalcanton et al. 2012) and the Legacy Extra-Galactic UV Survey (LEGUS; Calzetti et al. 2015). The former focuses on M31 to extreme depth, providing an extensive and very complete cluster catalog, but for only a single galaxy. The latter targets 50 nearby galaxies chosen from across the star-forming portion of the Hubble sequence, providing data on a much greater range of galactic environments, but with less depth. A preliminary cluster catalog for PHAT is now available (Johnson et al. 2012), and catalogs for the LEGUS galaxies will be released in the future (A. Adamo et al., 2015, in preparation).

The theoretical challenge when working with these extended cluster samples is that deriving masses and other properties for small clusters is non-trivial because they do not fully sample the stellar initial mass function (IMF) (e.g., Cerviño & Valls-Gabaud 2003; Cerviño & Luridiana 2004, 2006; Maíz Apellániz 2009). Traditional methods of determining star cluster properties by isochrone fitting therefore begin to fail, because at such small masses the relationship between clusters’ photometric properties and their mass, age, and other physical properties is non-deterministic. One approach to dealing with this problem is to resolve at least the most massive members of the cluster and determine their properties star-by-star using a color–magnitude diagram (CMD; Beerman et al. 2012). However, for young open clusters where the stellar surface density is high and confusion is significant, this requires extreme resolution and depth in the observations, and thus is not yet practical as a method for obtaining large cluster samples across a wide range of galactic environments.\(^{17}\)

The alternative approach is to develop new statistical techniques that can cope with partial sampling, and this has motivated the development of three major stochastic stellar population synthesis and analysis codes: MASClean (Popescu & Hanson 2009, 2010a, 2010b), a stochastic version of Pegase (Fouesneau & Lançon 2010; Fouesneau et al. 2012, 2014), and slug (da Silva et al. 2012, 2014; Krumholz et al. 2015). While these codes use slightly different statistical techniques, they all attempt to solve essentially the same problem: given a set of observed photometric properties for a star cluster, what should we infer about the probability distribution for its mass, age, extinction, or other physical properties?

In this paper we use slug, the Stochastically Lighting Up Galaxies code, and its post-processing tool for analysis of star cluster properties, cluster slug, to analyze an initial sample of clusters from the LEGUS (Calzetti et al. 2015). The clusters were identified in three HST WFC3 pointings (field of view $2\arcmin \times 2\arcmin$). Two pointings span the radius of the inner disk of NGC 7793, and include its nucleus, while the third pointing is in NGC 628, just inside of the eastern inner disk edge. Both NGC 628 and NGC 7793 are well-studied late-type spiral galaxies, and are relatively isolated relative to other massive systems. However, they provide complementary views of star cluster populations. Whereas NGC 7793 is in the nearby Sculptor group at 3.6 Mpc, and is inclined, NGC 628 is nearly face-on and is more distant, at 9 Mpc (Tully et al. 2009). Their star formation rates and stellar masses also differ by a factor of $\sim 3$, and are $0.7 M_\odot$ yr$^{-1}$ and $4 \times 10^7 M_\odot$ for NGC 7793, and $2.0 M_\odot$ yr$^{-1}$ and $1.6 \times 10^8 M_\odot$ for NGC 628 (Lee et al. 2009; Cook et al. 2014). Thus NGC 628 and 7793 represent examples toward the extremes of the full LEGUS sample of late-type spirals in terms of distance, orientation, mass, and star formation rate. This makes them a useful testbed for the performance of our analysis, one that will be extended to the full LEGUS sample in future work.

We have two goals in this paper. First, we wish to understand the performance of the stochastic stellar population analysis, including how it compares to traditional fitting methods, and to understand any systematic errors that might appear in the results due to choices of stellar tracks, stellar IMF, extinction law, assumed priors, or similar issues. Analysis of this sort has previously been performed by Popescu et al. (2012) for the Large Magellanic Cloud, by Fouesneau et al. (2012, 2014) for M31, and by de Meulenaer et al. (2013, 2014, 2015) for M31 and M33, but studies of systems have mostly been limited to questions of isolated metallicity, evolutionary tracks, and extinction law. To our knowledge no previous work has covered the possible systematics thoroughly in the context of a single framework, and none at all have explored the issue of priors. Second, we wish to describe a portion of the LEGUS data pipeline, which will be used to produce star cluster catalogs based on cluster slug modeling for the entire LEGUS sample. While this paper focuses on three fields in two galaxies for early analysis, the methods we develop here will be applied to the entire LEGUS sample. Ultimately, we will provide a catalog of homogeneously observed and analyzed star clusters, whose properties have been derived with a fully stochastic treatment of stellar population synthesis, for 50 galaxies that sample across the star-forming part of the Hubble sequence. This data set will provide tremendous insight into the extent to which star cluster formation and evolution depends on galactic environment.

The plan for the remainder of this paper is as follows. In Section 2 we describe the method we use to derive the physical properties of clusters, both stochastically and, for comparison, using a deterministic method. Section 3 presents the results of this analysis for individual clusters, and discusses general trends and possible systematics. In this section we also investigate how the stochastic and deterministic analyses compare. Section 4 extends this analysis to the properties of the star cluster population as a whole. Finally, Section 5 summarizes our conclusions.

2. METHODOLOGY

2.1. The Input Photometric Catalog

A description of the steps required to produce final cluster catalogs of the LEGUS targets can be found in Calzetti et al. (2015), and in A. Adamo et al. (2015, in preparation) we will present the custom pipelines we have developed to analyze the whole sample in a homogeneous fashion. We summarize here

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\(^{17}\)To our knowledge the largest-scale application of this technique to appear in the literature to date is in Fouesneau et al. (2014), who use PHAT photometry to obtain CMD-based ages for 100 clusters in M51. However, this paper presents the CMD-based results only for the purposes of comparison with those based on integrated photometry, and does not contain the CMD analysis itself, which is deferred to an as-yet unpublished paper.
the main aspects of this process which led to the final cluster catalogs of the three pointings in two galaxies used in this work, i.e., NGC 628 East (hereafter, NGC 628e), NGC 7793 West (NGC 7793w), and NGC 7793 East (NGC 7793e). We create our catalogs by using the SExtractor algorithm (Bertin & Arnouts 1996) on white-light images produced with the five standard HST LEGUS UV–UBVI bands (WFC3 F275W and F336W, WFC3 or ACS F438W or F435W, F555W, and F814W). The exact filters and exposure times used in each pointing are listed in Table 1. The cluster candidate catalogs contain only sources which satisfy the following criteria: (1) the V band concentration index (CI) must be greater than the stellar CI peak (the CI reference value slightly changes as function of galactic distance and HST camera); (2) the cluster candidate must be detected in two contiguous bands with a signal-to-noise ratio higher than 3.

The multi-band photometry of the cluster candidates of the three galactic fields used in this work has been derived with a fixed photometric aperture. We used an aperture radius of 4 and 5 px for NGC 628 and NGC 7793, respectively. The aperture radius is determined to include at least 50% of the flux of a median built growth curve of several isolated clusters selected in the field of each galaxy. The value will change mainly as a function of galactic distance. The sky annulus 1 pixel wide is located at 7 pixel from the center of the source. Averaged aperture corrections in all the filters have been estimated using manually selected isolated clusters in each galactic field. To take into account uncertainties produced by using a fixed aperture correction, we have added in quadrature the standard deviations of the aperture correction and the raw photometric error.

To remove contaminants from these automatically created catalogs we have visually inspected a subsample of cluster candidates in each galactic field with detections in at least four filters and absolute magnitudes brighter than $M_V = -6$ mag. Each inspected source has then been assigned a morphology quality flag. Class 1 objects are compact and centrally concentrated clusters; class 2 systems are clusters which show slightly elongated density profiles; class 3 are systems which show asymmetric profiles and multiple peaks, suggesting an association of stars; class 4 objects are single stars or artifacts, or any other spurious detection which can be excluded from the cluster catalog. In this work we will only use high-confidence, visually confirmed cluster catalogs clusters with photometry in at least 4 filters and flagged as class 1–3. However, as a service for readers who might wish to make their own classifications and produce their own samples, in the electronic edition of the paper we provide a separate machine-readable table containing the results of a cluster slug analysis of the clusters we omit here. We also note that, given their morphology, one might question whether class 3 objects should be called “clusters.” For the purposes of this analysis we treat them as such, but the fact that we have included them should be kept in mind when examining our results on the cluster population. Excluding them would produce a population that includes somewhat fewer young objects, since we find slightly younger than average ages for class 3 objects. However, since the focus of this work is on the method of analysis rather than the details of the population, we defer further discussion of this topic to future work.

Before proceeding, we make two additional notes regarding the catalog. First, we exclude the nuclear star cluster of NGC 7793 from all the analysis presented below. Although this cluster is flagged to be in the high-confidence catalog, it is almost certainly not a simple stellar population, and we should not analyze its properties as if it were. Second, due to partial overlap of the NGC 7793e and NGC 7793w pointings, some clusters appear twice in the catalog, once for each field. We have retained the two separate entries in the analysis below, since they are observed with slightly different filters and thus do not produce completely identical results. However, when presenting population statistics in Section 4, we include only the version of the cluster that appears in our NGC 7793e catalog.

The final, high-confidence cluster catalog we use for this paper contains 621 distinct clusters. Counting the duplicates that appear in the overlapping fields for NGC 7793 separately, the total rises to 645.

### 2.2. Stochastic Models with cluster_slug

Our stochastic analysis makes use of the cluster_slug software package that is part of the slug (Stochastically Lighting Up Galaxies) software suite (da Silva et al. 2012, 2014; Krumholz et al. 2015). The Bayesian analysis performed by cluster_slug takes as input a set of absolute magnitudes $M_F$ in some set of filters, with corresponding errors $\Delta M_F$. It returns as output the posterior probability distribution for the cluster mass $M$, age $T$, and extinction $A_V$. Details of how this operation is performed can be found in the papers referenced above, and both the slug suite and the full software pipeline used in this paper are available at http://www.sluagsps.com. Here we limit our discussion to the choice of two key inputs for the analysis: the libraries of simulated star clusters on which cluster_slug operates and the choice of prior probabilities that enter any Bayesian analysis. For a third input, the kernel density estimation bandwidth, we use $h = 0.1$ dex in the physical variables and $h = 0.1$ mag in the photometric ones.

The cluster_slug code requires a library of simulated star clusters to act as a “training set.” We produce these using the slug stochastic stellar population synthesis code. To generate a library, we must specify the stellar evolution tracks, metallicity, extinction curve, stellar IMF, and the fraction of the ionizing light that is reprocessed into nebular emission within the observational aperture. One of the goals of this paper is to study how these choices affect the deduced cluster properties, and for this reason we generate a range of libraries, summarized in Table 2. The parameters we vary in these libraries are as follows: for tracks we consider both those from the Padova

| Field      | Filter  | $t_{exp}$ (s) | Filter  | $t_{exp}$ (s) | Filter  | $t_{exp}$ (s) | Filter  | $t_{exp}$ (s) | Filter  | $t_{exp}$ (s) |
|------------|---------|---------------|---------|---------------|---------|---------------|---------|---------------|---------|---------------|
| NGC 628e   | WFC3    | 2361          | WFC3 F336W | 1119          | ACS F435W | 4720          | WFC3 F555W | 965          | ACS F814W | 1560          |
| NGC 7793w  | WFC3    | 2349          | WFC3 F336W | 1011          | WFC3 F435W | 947           | WFC3 F555W | 680          | WFC3 F814W | 430           |
| NGC 7793e  | WFC3    | 2349          | WFC3 F336W | 1011          | WFC3 F435W | 947           | ACS F555W  | 1125         | ACS F814W | 971           |

**Table 1**

List of Filters and Exposure Times

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The Astrophysical Journal, 812:147 (23pp), 2015 October 20

Krumholz et al.
group including TP-AGB stars (Vassiliadis & Wood 1993; Girardi et al. 2000; Vázquez & Leitherer 2005) and from the Geneva group (Ekström et al. 2012). For metallicities we consider \( Z = 0.004, 0.008, 0.020 \) for the Padova tracks, and \( Z = 0.014 \) for the Geneva tracks.\(^{18}\) We consider extinction curves appropriate to both the Milky Way (Landini et al. 1984; Fitzpatrick 1999) and to a starburst (Calzetti et al. 2000), and both Kroupa (2001) and Chabrier (2005) IMFs. Finally, for nebular emission we consider both \( \phi = 0.5 \) (50% of ionizing photons available to produce nebular emission inside the aperture) and \( \phi = 0 \) (no nebular emission). Figure 1 provides examples of how some of these choices affect the predicted variation of magnitude versus age in select filters. All the examples shown are for a \( 10^5 M_\odot \) cluster with a fully sampled Kroupa IMF, and no extinction. One can clearly see in the Figure 1 the differences in predicted luminosity due to the treatment of TP-AGB stars at older ages, and due to the treatment of nebular emission at younger ages.

Our fiducial library, which we use for all purposes unless specified otherwise, is pad_020_kroupa_MW. This uses Padova isochrones, metallicity \( Z = 0.02 \), a Kroupa IMF, a Milky Way extinction curve, and 50% of ionizing photons converted to nebular emission. We adopt this as our fiducial model because it appears to be the closest match to what we know of the target galaxies. Observed IMFs in local spiral galaxies are consistent with a universal function consistent with the Kroupa (2001) parameterization (see the recent reviews by Kroupolz 2014 and Offner et al. 2014). Both NGC 628 and NGC 7793 have mean metallicities close to Solar, and for the typical \( \approx 0.03 \) dex kpc\(^{-1} \) metallicity gradients observed in the inner disks of local spirals (e.g., Pilyugin et al. 2004), the \( <10 \) kpc-wide regions spanned by our observed fields should have end-to-end metallicity differences of only a few tenths of a dex. On average we observe that, for the deterministic models at least (see below), clusters within face-on Solar metallicity spirals are better fitted by a Milky Way extinction curve (Cardelli et al. 1989) than with the alternatives. Finally, we use the Padova rather than Geneva evolutionary tracks because the majority of our clusters are older than \( \approx 10 \) Myr, and in this age range the superior treatment of AGB stars in the Padova models is advantageous. All these fiducial choices are discussed further in A. Adamo et al. (2015, in preparation).

In addition to these choices, for each library we must specify how the masses, ages, and extinctions of the simple stellar populations are to be sampled. For all the libraries used in this paper, we choose a distribution that varies with cluster mass and age as

\[
p_{\text{hh}}(x) = p_M(\log M)p_T(\log T)p_{A_V}(A_V)
\]

where

\[
p_M(\log M) = \begin{cases} 1, & 2 < \log M < 4 \\ 10^{-\log M - 4}, & 4 \leq \log M < 8 \end{cases}
\]

\[
p_T(\log T) = \begin{cases} 1, & 5 < \log T < 8 \\ 10^{-\log T - 8}, & 8 < \log T < \log T_{\text{max}} \end{cases}
\]

\[
p_{A_V}(A_V) = \begin{cases} 1, & 0 < A_V < 3 \end{cases}
\]

In the above expressions \( M \) is in units of \( M_\odot \), \( T \) is in units of years, and \( A_V \) is in mag. We use \( T_{\text{max}} = 15 \) Gyr for models using the Padova tracks, and \( T_{\text{max}} = 1 \) Gyr for models using Padova Kroupa IMF.

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\(^{18}\) In principle it would be best to treat metallicity as a continuous variable to be fit, as we will fit cluster mass and age. However, tracks including TP-AGB stars that are sufficiently finely sampled in metallicity to make such a treatment possible did not become available from the Padova group until late 2014, and have not yet been incorporated into slug. No correspondingly fine tracks are available from the Geneva group yet. As a result, we will treat metallicity as fixed. As discussed below, this is not a significant limitation for our target galaxies, because the expected metallicity variation within the galaxy is relatively small.
the Geneva tracks. This sampling distribution places most realizations at small masses and ages, where stochastic variation is largest, and uses fewer realizations at higher masses and older ages where they are smaller. We generate $10^7$ realizations for each library.

The final ingredient needed for a cluster slug calculation is a prior probability distribution for the physical properties, which we use to weight the simulations in the library. One can demonstrate that the choice of priors will be important in at least some regimes via a simple thought experiment. Consider what might seem like a natural prior, a distribution that is flat in the log of the age. Stars’ luminosities and colors do not change significantly until their ages reach $\sim1$ Myr. Thus all star clusters younger than this age look identical, and the shape of the posterior probability distribution at ages below $\sim1$ Myr must therefore be the same as the shape of the prior probability distribution. For a prior that is flat in log age, we would therefore conclude that age ranges of $\log(T/\text{year}) = 3.5 - 4$ and $\log(T/\text{year}) = 5.5 - 6$ are equally likely. This seems unlikely to be correct: even if it were completely unbound, a stellar population formed $<1$ Myr ago would still be close enough together to be recorded as a cluster in the LEGUS catalog, so every cluster that reaches an age of $\log(T/\text{year}) = 3.5 - 4$ must eventually reach an age of $\log(T/\text{year}) = 5.5 - 6$. However, this cluster will reside in the interval $\log(T/\text{year}) = 3.5 - 4$ for a mere 6800 years, while it will spend 680,000 years in the interval $\log(T/\text{year}) = 5.5 - 6$. Clearly there should be 100 times as many clusters in the latter bin as in the former, making the latter 100 times as likely. Even worse, suppose further that we have a cluster whose colors admit ages older as well as younger than $\sim1$ Myr. In this case the weight that we would end up assigning to the larger ages would depend strongly on whether our model grid contained a youngest age of, say, $10^3$ years, $10^4$ years, $10^5$ years, or $10^6$ years, because the amount of probability “phase space” allowed at young ages would depend on this choice. Clearly some attention to priors is required to avoid outcomes such as this.

For ages younger than a few Myr, the proper prior is almost certainly flat in age. This is because, even if a newly formed collection of stars has negligible gravitational binding energy, at the typical velocity dispersions of a few km s$^{-1}$ found in young clusters, it will disperse over a region no more than $\sim10$ pc in size over this time. In an extragalactic survey such as LEGUS, it would still be detected as a cluster. Thus cluster dispersal is irrelevant at such young ages. Similarly, as noted above, stellar evolution effects are also negligible. Given the absence of any physical mechanism to bias the age distribution, all ages are equally likely, suggesting that the proper prior is flat in age.

For ages greater than $\sim10^{6.5}$ years, the proper prior is significantly less certain, as there is an extensive debate in the literature on this topic. Some authors report distributions $dN/dT \propto T^{-1}$ (i.e., flat in log age; e.g., Fall et al. 2005, 2009; Chandar et al. 2010a, 2010b, 2011; Fall & Chandar 2012; Fousneau et al. 2012; Popescu et al. 2012), while others (mostly focusing on ages > 10 Myr) find $dN/dT \propto T^9$ (i.e., flat in age; e.g., Boutloukos & Lamers 2003; de Grijs et al. 2003a; Gieles et al. 2007; Bastian et al. 2011, 2012a, 2012b; Silva-Villa et al. 2014); some authors also report intermediate results, with the index of the age distribution changing as a function of age (e.g., Fousneau et al. 2014). The results appear to depend in part on the criteria one uses to define what is a cluster and thus should be included in the cluster catalog; see Krumholz (2014) for a recent review. The situation is further complicated by the fact that the prior we want to use is not the intrinsic distribution of star cluster ages, but the convolution of that with our selection function. Properly modeling this would require much greater understanding of our observational completeness and biases than we now possess. Given these complexities, as a fiducial prior we choose a compromise, $dN/dT \propto T^{-0.5}$, which is intermediate between flat in age ($dN/dT \propto$ constant) and flat in log age ($dN/d\log T \propto$ constant, so that $dN/dT \propto T^{-1}$). However, we will explore the effects of varying this choice below. The maximum possible age will be set by the maximum age in our library.

For the prior on the mass distribution, we note that observations of young star clusters consistently find $dN/dM \propto M^{-2}$ (e.g., Williams & McKee 1997; Zhang et al. 1999; Larson 2002; Bik et al. 2003; de Grijs et al. 2003b; Goddard et al. 2010; Bastian et al. 2011; Fall & Chandar 2012; Fousneau et al. 2012), with relatively little variation. These results are mostly derived at masses high enough that stochasticity is unimportant, and thus it seems reasonable simply to extrapolate them into our regime. We therefore adopt $p(M) \propto M^{-1}$ (corresponding to $dN/dM \propto M^{-2}$) as our fiducial prior on the mass distribution. Two caveats are worth mentioning, though. First, as with the age distribution, the truly correct prior to use would be the convolution of this function with our observational selection, which unfortunately is not fully characterized as yet. Second, to avoid a logarithmic divergence, we must truncate the prior distribution at both the high and low mass ends; we implicitly choose the truncation values via the range of masses we include in our library. Fortunately the choices here have relatively little effect, because our library covers masses from $10^2$ to $10^8 M_\odot$, which spans the entire plausible mass range: the lower limit of $100 M_\odot$ is almost certainly smaller than the smallest cluster we have any hope of detecting, and, given the size of our sample, the upper limit is much more massive than the most massive cluster we would expect to find even if the underlying cluster mass function continued as $dN/dM \propto M^{-2}$ to even higher masses.

Finally, for the $A_V$ distribution, observations indicate that $A_V$ values tend to be at most $\sim0.5$ at ages above $\sim10$ Myr (e.g., Whitmore & Zhang 2002; Mengel et al. 2005; Whitmore et al. 2011; Bastian et al. 2014; Hollyhead et al. 2015). At younger ages, Fousneau et al. (2012, 2014) report a fairly broad distribution of $A_V$. In principle we could use a prior that combines age and $A_V$ to disfavor a combination of old age and high-$A_V$. However, the functional form of this constraint is not well-characterized, and we wish to keep our priors on age and $A_V$ independent to ease the analysis of how our results depend on a choice of priors. For this reason we choose to adopt a flat distribution in $A_V$ as a prior.

In summary, our fiducial prior distribution follows $p_{\text{prior}}(x) \propto M^{-1} T^{-0.5}$, with no dependence on $A_V$. However, in Section 3.3 we will check the dependence of our results upon this choice.

2.3. Deterministic Models and Traditional $\chi^2$-fitting Method

The standard photometric analysis of the LEGUS cluster catalogs is performed with Yggdrasil\textsuperscript{19} (Zackrisson

\textsuperscript{19} http://ttt.astro.su.se/projects/yggdrasil/yggdrasil.html)
et al. 2011) deterministic models and a traditional \( \chi^2 \)-fitting procedure (e.g., see Calzetti et al. 2015). Yggdrasil is a spectral synthesis code that can provide spectral and integrated flux tables for a large set of stellar and nebular parameters. In Yggdrasil, the spectrum produced by the stellar population at each age step is then used as input for a calculation of the nebular emission using cloudy (Ferland et al. 2013) to ensure a self-consistent treatment.

For the LEGUS cluster analysis in this paper, we have carefully matched the physical models used in Yggdrasil to those of the fiducial cluster_slug library. Thus for our Yggdrasil models we assume that clusters are a simple stellar population, and we compute the stellar spectrum using Starburst99 (Leitherer et al. 1999; Vázquez & Leitherer 2005) atmosphere models. We use the Kroupa (2001) IMF, the same Padova AGB stellar tracks used in the cluster_slug models, and a metallicity \( Z = 0.02 \). We calculate nebular emission assuming that (1) the metallicity of gas and stars is the same, (2) the average gas density is 100 cm\(^{-3} \), (3) the ionizing flux is normalized by setting it to that of a 10\(^6\) \( M_\odot \) stellar population, and (4) the covering factor set to 50%, meaning that 50% of the ionizing photons are assumed to be absorbed by hydrogen atoms within the observational aperture. To model extinction, we adopt a Milky Way extinction curve, and consider differential extinctions \( E (B-V) \) between 0 and 1.5 mag with steps of 0.01 mag. We apply extinction to the model spectra before convolving them with the filter throughput. Thus in every respect except nebular emission, which we discuss in Section 2.3.1, our Yggdrasil models are completely matched to our fiducial cluster_slug one.

The \( \chi^2 \)-fitting algorithm used to derive cluster physical properties has been presented and tested in Adamo et al. (2010). The algorithm compares the observed photometry to the library of Yggdrasil models, and returns the library cluster age, extinction, and mass that deviates least from the observations. A quality of the fit value is estimated from the reduced \( \chi^2 \), i.e., \( Q \)-parameter. Good fits have \( Q \) values close to 1, while values below 0.1 suggest poor fits to the observed cluster SEDs. The \( \chi^2 \)-fitting algorithm has also been extended by Adamo et al. (2012) to produce uncertainties in the derived cluster physical properties contained within the 68% confidence limits of the solution parameter space (Calzetti et al. 2015, their Figure 9).

### 2.3.1. A Caveat Regarding Nebular Emission

Before presenting results, we pause to add a caveat to our comparison of the results obtained with Yggdrasil and cluster_slug. We have tried to ensure that the underlying physical models are as similar as possible in the two codes, so that any differences between them arise purely from the fact that cluster_slug includes stochasticity and returns a full posterior PDF, while our models using Yggdrasil are non-stochastic, and our fits to them use a traditional \( \chi^2 \) method. The one quantity we have not been able to match precisely is the treatment of nebular emission.

To good approximation, an \( \text{H}\alpha \) region produces a nebular luminosity per ionizing photon injected that is a function of the gas metallicity and ionization parameter. While the Yggdrasil and cluster_slug models are matched in metallicity, they are not precisely matched in ionization parameter. As noted above, the nebular contribution in the Yggdrasil models we use in this paper has been derived by taking spectra computed for a 10\(^6\) \( M_\odot \) stellar population illuminating an \( \text{H}\alpha \) region with a constant density of 10\(^2\) cm\(^{-3} \) and a 50% covering fraction. (See Zackrisson et al. 2011 for further details.) In the Yggdrasil models, the inner radius of the \( \text{H}\alpha \) region scales with total luminosity to the 1/2 power, so the ionization parameter at the illuminated face of the \( \text{H}\alpha \) region is kept roughly constant—roughly rather than precisely because the ratio of ionizing to total luminosity is not constant as a stellar population ages. However, because the volume of ionized gas changes as the luminosity does, and the ionizing photon flux changes as the photons propagate through the nebula, this prescription causes the volume-averaged ionization parameter to vary significantly with the mass and age of the stellar population. The mean is close to log \( U = -2.5 \), but there is a non-negligible spread.

Slugg’s treatment of nebular emission also assumes constant density within an \( \text{H}\alpha \) region, and we have chosen the same 50% covering fraction as in the Yggdrasil models. However, in its simplest form, slug assumes that the volume-averaged ionization parameter has a constant value log \( U = -3 \), while the ionization parameter at the inner edge is not held constant.20 While slug also supports an option to pass the spectra through cloudy, and thus in principle we could fully emulate Yggdrasil’s assumptions, it is not computationally feasible to run cloudy on all, or even a substantial subset, of the nearly 10\(^8\) models in our libraries. Thus we are left with an imperfect match in ionization parameter between the cluster_slug and Yggdrasil models; the Yggdrasil models assume that \( \text{H}\alpha \) regions form a sequence where the ionization parameter at the illuminated face is fixed but the volume-averaged ionization parameter is not, while the cluster_slug models fix the volume-averaged but not the illuminated face ionization parameters.

How does this difference in ionization parameter translate into a difference in photometry, which is what we ultimately care about? Obviously the difference will be negligible in stellar populations older than \( \sim 4 \) Myr, when the ionizing luminosity drops precipitously. Even in younger populations, the nebular contribution to the total luminosity is only \( \sim 10\% \) for F275W and F438W filters, and thus a difference in how this emission is treated is again unimportant. The nebular contribution is most significant for F336W, F555W, and F814W, where it can reach \( \sim 60\% - 70\% \) of the total (e.g., Reines et al. 2010).

In the case of F555W, the nebular contribution is dominated by the [O ii] \( \lambda \lambda 4959, 5007 \) and H\beta lines, with additional contributions from bound-free and two-photon continuum emission, and from the H\alpha and [N ii] \( \lambda \lambda 6584 \) lines (which are intrinsically bright but lie near the edge of the filter response curve, and will be shifted out of the filter entirely for even modest redshifts). Over the observationally plausible range of ionization parameters (roughly log \( U = -3 \) to \( -2 - \ldots \), e.g., see the compilation in Verdolini et al. 2013), both Balmer line and bound-free production per ionizing photon injected vary by a factor of \( \sim 2 \). The [O ii] emission, as well as the other metal lines, are sensitive to both the temperature and the ionization state of the nebula. Consequently, they can easily vary by an order of magnitude as the ionization parameter changes (e.g., Kewley et al. 2001; Verdolini et al. 2013; Yeh et al. 2013), and

20 Full details of the slug method, and the computational motivations for choosing a constant ionization parameter, are given in Krumholz et al. (2015).
at the highest ionization parameters they can be a factor of several brighter than the H β line. For F814W, the nebular contribution is dominated by free–free and bound–free continuum, with an a sub-dominant contribution of Paschen lines and [S iii] λ9069, 9532. The two continuum sources and the Paschen lines are relatively insensitive to the ionization parameter, varying by less than a factor of 2. In contrast, the [S iii] lines, while they are subdominant, have the largest potential variation. Their strength can change by at least an order of magnitude over the plausible range of ionization parameters. Finally, the nebular contribution F336W is strongly dominated by the bound-free and two-photon emission. The former varies by less than a factor of 2 over the plausible ionization parameter range, while the latter stays almost entirely constant until the density reaches \( \sim 10^5 \text{ cm}^{-3} \), at which point the two-photon process is collisionally quenched.

Taken together, these factors suggest that differences in the assumed ionization parameter between Yggdrasil and cluster_slug will lead to at most factor of \( \sim 3 \) differences in F336W, F555W, and F814W luminosity in young stellar populations. These variations are not trivial, but we shall see below that they constitute a fairly minor uncertainty in comparison to those associated with degeneracies between age and extinction, or variations due to stochasticity. However, as a final caution, we note that the prescriptions for nebular emission in both cluster_slug and Yggdrasil are only rough approximations to the real complexity of H ii regions. For example, both cluster_slug and Yggdrasil assume a fixed, age-independent covering fraction. However, in reality H ii regions expand with time while the observational aperture remains fixed, so the covering fraction should drop with age (cf. Whitmore et al. 2011), at a rate that depends on both the cluster’s luminosity and the density of its surrounding, both of which affect the rate at which its H ii regions expands. Thus age estimates below \( \sim 5 \text{ Myr} \) should be regarded with extreme caution regardless of the fitting method.

3. RESULTS FOR INDIVIDUAL CLUSTERS

Throughout this section we will refer to the results of an analysis of the cluster populations of NGC 7793e, NGC 7793w, and NGC 628e performed using cluster_slug. For the convenience of other authors who wish to use our results without having to rerun the full analysis, we summarize the cluster_slug output in Table 3. A full machine-readable version of this table is included in the electronic edition of this paper, along with an analogous table containing the same results for the objects we have classified as unlikely to be genuine star clusters. We provide the latter as a service for those who might wish to make their own classifications.

3.1. Are the cluster_slug Libraries Consistent with the Observations?

As a first, most basic question, we examine the extent to which the libraries of stochastic models reproduce the colors and magnitudes found in the observed sample. The kernel density estimation method used by cluster_slug will return results even if the correspondence between the model libraries and the observations is very poor, but those results should be considered reliable only to the extent that the models and observations are in reasonable correspondence. To check whether this is the case, in Figure 2 we show the distribution of our cluster_slug model libraries in color–color space, as compared to the observed catalog. As the plot shows, the model libraries generally cover loci in color–color space very similar to the observations.

By itself, agreement in cuts in color–color space does not confirm that our libraries are a good analog to the observations. Even if the real and synthetic data show similar distributions in particular projections, they may be differently distributed in higher dimensions. Moreover, unlike in deterministic models like Yggdrasil, the total mass and absolute magnitude are not free parameters in the cluster_slug models. Because the amount of stochasticity depends on how well the mass distribution is sampled, the dispersion in color at fixed age is a function of the absolute magnitude. As a result, it is possible that our models cover the same range as the data in color–color space, but that they might not in color–magnitude space.

To perform a more quantitative comparison of the observed and synthetic catalogs, we therefore examine the distribution of distances in photometric space between the observed star clusters and their closest analogs in our synthetic catalogs. We define the absolute and normalized photometric distances between an observed star cluster and the \( i \)th member of a cluster_slug library by

\[
D = \sqrt{\frac{1}{N} \sum_j \left( M_{F,i,\text{obs}} - M_{F,j} \right)^2 },
\]

\[
D_{\text{norm}} = \sqrt{\frac{1}{N} \sum_j \left( \frac{M_{F,i,\text{obs}} - M_{F,j}}{\Delta M_{F,i,\text{obs}}} \right)^2 },
\]

where \( N \) is the number of filters, \( M_{F,\text{obs}} \) is the magnitude of the observed cluster in filter \( F \), \( \Delta M_{F,i,\text{obs}} \) is the observational error on this value, and \( M_{F,j} \) is the magnitude of the \( j \)th library cluster in filter \( F \). The exact filters used in this comparison are those given in Table 1, so generally \( N = 5 \); however, there are a very small number of clusters which lie outside the image in one of the filters, and for these \( N = 4 \).

For each observed cluster, we define \( D_5 \) and \( D_{\text{norm}}^5 \) as the absolute and normalized distances to the 5th nearest neighbor in one of the cluster_slug libraries. (Note that the 5th nearest neighbor is not necessarily the same simulated cluster for the absolute and normalized distances.) Figure 3 shows the cumulative distribution of \( D_5 \) and \( D_{\text{norm}}^5 \) for each of our sample galaxies compared to each of our model catalogs; plots for the \( n \)th nearest neighbor, with \( n \sim 1–20 \), are qualitatively similar. Clearly nearly all of our observed star clusters have close analogs in our synthetic catalogs. For our fiducial model library, pad_020_kroupa_MW, we find that 71% of observed clusters have at least 5 matches in our library within the \( 1\sigma \) photometric errors (i.e., 71% have \( D_5^\text{norm} < 1 \)), 95% have at least 5 matches within the \( 2\sigma \) photometric errors, and 99% have 5 matches within the \( 3\sigma \) errors. The figures are comparable for the other model libraries, and the differences between them are quite minor. The normalized distances are generally smallest for NGC 0628e and largest for NGC 7793e, but this is more a reflection of the size of the photometric errors in those fields than of any intrinsic differences between the cluster populations.

While the nearness of our library points to the observations is encouraging, we should inquire a bit further about what the distribution of \( n \)th nearest neighbor distances should look like
### Table 3
Summary Marginal Posterior PDFs

| ID           | Mode Class<sup>a</sup> | Mean Class<sup>a</sup> | log($M/M_{\odot}$) Percentiles | log($T$/year) Percentiles | $A_V$ Percentiles | $D_5$ (mag) | $D_5^{\text{norm}}$ |
|--------------|------------------------|------------------------|-------------------------------|---------------------------|------------------|-------------|------------------|
|              |                        |                        | 16   | 25   | 50   | 75   | 84   | 16   | 25   | 50   | 75   | 84   | 16   | 25   | 50   | 75   | 84   | 16   | 25   | 50   | 75   | 84   |
| pad_020_kroupa_MW, $\beta = -2.0$, $\gamma = -0.0$ | 3 | 1.00 | 1.00 | 4.33 | 4.38 | 4.48 | 4.58 | 4.63 | 8.10 | 8.18 | 8.30 | 8.42 | 8.46 |
|               |                        |                        | 4   | 1.00 | 1.00 | 3.79 | 3.84 | 3.99 | 4.09 | 4.19 | 6.14 | 6.18 | 6.34 | 6.51 | 6.71 |
|               |                        |                        | 5   | 2.00 | 1.67 | 2.50 | 2.65 | 3.00 | 3.34 | 3.49 | 6.22 | 6.30 | 6.43 | 6.59 | 6.71 |
| pad_020_kroupa_MW, $\beta = -2.0$, $\gamma = -0.5$ | 3 | 1.00 | 1.00 | 3.44 | 4.29 | 4.43 | 4.53 | 4.58 | 7.00 | 8.02 | 8.26 | 8.38 | 8.42 |
|               |                        |                        | 4   | 1.00 | 1.00 | 3.79 | 3.84 | 3.94 | 4.09 | 4.14 | 6.10 | 6.18 | 6.30 | 6.47 | 6.55 |
|               |                        |                        | 5   | 2.00 | 1.67 | 2.50 | 2.65 | 2.95 | 3.25 | 3.39 | 6.22 | 6.30 | 6.43 | 6.55 | 6.63 |
| pad_020_kroupa_MW, $\beta = -2.0$, $\gamma = -1.0$ | 3 | 1.00 | 1.00 | 2.85 | 3.05 | 3.54 | 4.48 | 4.53 | 6.75 | 6.79 | 7.04 | 8.26 | 8.34 |
|               |                        |                        | 4   | 1.00 | 1.00 | 3.79 | 3.84 | 3.94 | 4.04 | 4.14 | 6.10 | 6.18 | 6.30 | 6.43 | 6.51 |
|               |                        |                        | 5   | 2.00 | 1.67 | 2.45 | 2.60 | 2.95 | 3.20 | 3.34 | 6.22 | 6.26 | 6.39 | 6.51 | 6.59 |
| pad_020_kroupa_MW, $\beta = -2.0$, $\gamma = -0.0$ | 3 | 1.00 | 1.00 | 4.34 | 4.39 | 4.49 | 4.59 | 4.64 | 8.10 | 8.18 | 8.30 | 8.38 | 8.46 |
|               |                        |                        | 4   | 1.00 | 1.00 | 3.85 | 3.90 | 4.05 | 4.19 | 4.24 | 6.14 | 6.22 | 6.34 | 6.55 | 6.83 |
|               |                        |                        | 5   | 2.00 | 1.67 | 2.61 | 2.81 | 3.11 | 3.40 | 3.55 | 6.22 | 6.30 | 6.43 | 6.63 | 6.79 |
| pad_020_kroupa_SB, $\beta = -2.0$, $\gamma = -0.0$ | 3 | 1.00 | 1.00 | 3.34 | 3.44 | 3.59 | 3.74 | 3.79 | 7.65 | 7.73 | 7.89 | 8.06 | 8.14 |
|               |                        |                        | 4   | 2.00 | 1.00 | 3.39 | 3.49 | 3.64 | 3.74 | 3.79 | 7.49 | 7.57 | 7.73 | 7.85 | 7.93 |
|               |                        |                        | 5   | 3.00 | 1.67 | 2.21 | 2.30 | 2.55 | 2.85 | 2.95 | 5.49 | 5.57 | 5.77 | 5.98 | 6.06 |
| pad_020_kroupa_SB, $\beta = -2.0$, $\gamma = -0.5$ | 3 | 2.00 | 2.00 | 3.30 | 3.39 | 3.59 | 3.74 | 3.79 | 7.61 | 7.65 | 7.85 | 8.02 | 8.10 |
|               |                        |                        | 4   | 2.00 | 1.00 | 3.34 | 3.39 | 3.59 | 3.74 | 3.79 | 6.79 | 7.32 | 7.65 | 7.81 | 7.89 |
|               |                        |                        | 5   | 3.00 | 1.67 | 2.21 | 2.30 | 2.55 | 2.85 | 2.95 | 5.49 | 5.57 | 5.77 | 5.98 | 6.06 |
| pad_020_kroupa_SB, $\beta = -2.0$, $\gamma = -1.0$ | 3 | 2.00 | 2.00 | 3.25 | 3.34 | 3.54 | 3.69 | 3.74 | 7.53 | 7.61 | 7.77 | 7.93 | 8.02 |
|               |                        |                        | 4   | 2.00 | 1.67 | 2.50 | 2.75 | 3.44 | 3.64 | 3.74 | 6.59 | 6.67 | 7.28 | 7.69 | 7.81 |
|               |                        |                        | 5   | 3.00 | 1.67 | 2.21 | 2.30 | 2.55 | 2.85 | 2.95 | 5.49 | 5.57 | 5.77 | 5.98 | 6.06 |
| pad_020_kroupa_SB, $\beta = -2.0$, $\gamma = -0.0$ | 3 | 2.00 | 2.00 | 3.40 | 3.50 | 3.65 | 3.80 | 3.85 | 7.61 | 7.69 | 7.85 | 8.02 | 8.10 |
|               |                        |                        | 4   | 2.00 | 1.67 | 3.35 | 3.45 | 3.65 | 3.75 | 3.80 | 7.32 | 7.53 | 7.73 | 7.85 | 7.93 |
|               |                        |                        | 5   | 3.00 | 1.67 | 2.22 | 2.32 | 2.61 | 2.91 | 3.01 | 5.49 | 5.61 | 5.82 | 5.98 | 6.06 |

<sup>a</sup> Mode and mean of the classifications given by the visual classifiers; 0 = source was not visually classified (too faint); 1 = symmetric, compact cluster; 2 = concentrated object with some degree of asymmetry or color gradient; 3 = diffuse or multiple peak system, possibly spurious alignment; 4 = probable spurious detection (foreground/background source, single bright star, artifact.

Notes. The full table includes 645 visually confirmed clusters, of which 621 are unique and 24 are overlapping between NGC 7793e and NGC 7793w. It also includes 2326 objects visually classified as unlikely to be clusters (provided in separate files).

(This table is available in its entirety in machine-readable form.)
if our models are in fact a good fit. For a Poisson process, the cumulative distribution function of $n$th nearest neighbor distances $r$ for a point at position $x$ is given by

$$d_n(r|x) = 1 - \sum_{k=0}^{n-1} \frac{1}{k!} \lambda^k (r,x)^k e^{-\lambda(r,x)},$$

where $\lambda(r,x)$ is the expected number of points in a ball of radius $r$ centered on $x$. Intuitively, this is simply the statement that the distribution of $n$th nearest neighbor distances is 1 minus the probability that there are $n-1$ or fewer points within a ball of size $r$ around the point in question. Note that we are free to use any metric to measure $r$, so we can use the normalized photometric distance as well as the absolute one. Evaluating the expectation value $\lambda$ requires knowledge of the shape of the underlying probability distribution around the point of interest. In principle we could evaluate this for each of our points from the kernel density estimate for the PDF, but doing so would be quite computationally intensive, and for small distances would likely depend on our choice of
bandwidth parameter. Instead, we can qualitatively check whether our nearest neighbor distribution is consistent with our models being a good match to the data by making a few simplifying assumptions that allow us to evaluate $\lambda$ analytically.

Suppose that our library is large enough that we expect there to be $\ell$ simulations from the library within a normalized photometric distance $D_{\text{norm}} = 1$ of each observed point. Further suppose that the library of simulation points lies along an $M$-dimensional manifold within our five-dimensional photometric space; for example, if the simulations near a particular point mostly lie along a line in photometric space, then the number within a given photometric distance increases linearly with distance, and $M = 1$. If they are mostly along a plane then $M = 2$, and so forth. In this case we have

$$\lambda(r) = \ell r^M,$$

and we can evaluate the expected $n$th nearest neighbor distribution $d_n(r)$ directly. Figure 4 shows how the measured distribution of nearest neighbor distances for our fiducial model library, pad_020_kroupa_MW, and the various color lines show the expected distribution of 5th nearest neighbor distances, computed using Equations (7) and (8) for various values of the expected number of library models within the photometric error circle, $\ell$, and the dimensionality of the data, $M$, as indicated in the legend.

Figure 4. Cumulative distribution function of 5th nearest neighbor photometric distances normalized by the observed photometric errors, $D_{\text{norm}}$. The thick black line shows the distribution we measure summed over all our sample fields and using our fiducial model library, pad_020_kroupa_MW. The various color lines show the expected distribution of 5th nearest neighbor distances, computed using Equations (7) and (8) for various values of the expected number of library models within the photometric error circle, $\ell$, and the dimensionality of the data, $M$, as indicated in the legend.

3.2. Posterior PDFs of Star Cluster Mass, Age, and Extinction

Having verified that our synthetic libraries produce reasonable-matches to the data, we now focus on our fiducial library, pad_020_kroupa_MW, and our fiducial prior distribution, $p_{\text{prior}}(\log M, \log T, A_V) \propto M^{-0.5}$, leaving a discussion of the dependence of the results on these choices to the subsequent sections. We analyze the entire photometric catalog described in Section 2.1, and for each cluster we compute the marginal posterior probability of mass, age, and extinction on a grid of 128 points each, covering the full range of each of these values present in our synthetic library.

The computation is relatively fast—deriving each posterior PDF requires $\sim 1$ CPU-second per cluster for most high-confidence clusters in the catalog; those with the largest photometric error bars, or that are not well-fit by any models in our catalog (as is the case for many of the visually rejected candidates, which we have nonetheless analyzed for completeness) may take up to a few tens of CPU-seconds, since large error bars require that we search a larger volume of parameter space. Overall, we find that deriving posterior PDFs for all $\sim 3000$ candidate clusters in the full catalog using a single cluster_slug library and sets of priors requires a few hours using a multi-core workstation; performing a similar analysis for the $\sim 600$ high-confidence clusters requires tens of minutes.

Figures 5–7 show sample marginal posterior distributions for mass, age, and extinction for 15 clusters per field in our 3 sample fields. The clusters shown are chosen to be uniformly distributed in the 50th percentile estimates of their log mass, log age, and extinction. From the plots, we can see that in most cases the cluster_slug models identify a fairly narrow range of possible masses for each cluster, with a typical interquartile range of $\sim 0.2–0.3$ dex on the posterior probability. The distributions are for the most part unimodal. However, we can see that there are a few cases where the posterior mass distribution is broader or even bimodal, or where there is a tail of probability extending to very different masses, so that the 10th or 90th percentile whiskers extend very far beyond the 1st–3rd quartile range.

In comparison, the posterior probability distributions of age are somewhat broader and more likely to be bimodal. In the bimodal cases there is often one peak at a relatively old age, and another at an age of $\sim 10^6.5$ years. This is particularly true for NGC 628, which has the broadest photometric errors. The relative weighting of the two possible age fits, as we shall see below, is not independent of our choice of priors. In these cases the median age may not be a good representation of the actual age, because the median occurs near a local minimum of the PDF that lies partway between the two peaks. The posterior PDFs of extinction are also quite broad. In some cases they are bimodal, while in others they display a single peak but with an extended tail.

The nature of these bimodal fits is illustrated in Figure 8, which shows various 2D projections of the posterior PDFs for one example cluster from NGC 628e, which is typical of many of the bimodal fits we find. As the figure shows, the data are consistent with two “islands” of probability. The young island corresponds to an extinction $A_V \sim 0.2–1.5$, age $T \lesssim 3\text{–}10$ Myr, and mass $M \sim 3000 M_\odot$, while the old one is centered near $A_V \sim 0.1$, $T \sim 500$ Myr, $M \sim 3–5 \times 10^5 M_\odot$. The color and luminosity of this cluster can therefore be fit well by either a relatively massive, extinction-free, old cluster, or a
younger, somewhat less massive cluster that is red due to greater extinction.

The conclusions that many clusters when analyzed stochastically show multiple probability maxima is not new, and has been pointed out previously by Fouesneau et al. (2012, 2014), de Meulenaer et al. (2013, 2014, 2015), and Krumholz et al. (2015). Indeed, one could obtain such a result even with a deterministic method, provided that one used the full posterior PDF rather than an approximation to it such as a Gaussian centered on the local minimum of $\chi^2$. The primary reason is that there are a number of places in color space where star clusters with disparate physical properties are nonetheless very similar in color. The result is a likelihood function that is not well approximated by a uni-modal Gaussian.

### 3.3. Dependence on Choice of Priors

To what extent do the results for the posterior probabilities depend on the choice of prior probability distribution? As noted above, the priors certainly matter for ages young enough that the color provides few constraints, but the priors may also matter in other parts of parameter space as well. To answer this question, we continue to use our fiducial library, but now consider prior probability distributions of the form

$$p_{\text{prior}}(x) \propto \begin{cases} M^{\beta+1} T, & \log(T/\text{year}) < 6.5 \\ M^{\beta+1+\gamma} \log(T/\text{year})^{+1}, & \log(T/\text{year}) \geq 6.5 \end{cases}$$

with $\beta = -1$ or $-2$, and $\gamma = 0, -0.5$, or $-1$; the combination $\beta = -2, \gamma = -0.5$ is our fiducial choice. That is, we consider cases where we take the prior on mass to be either flat in log mass ($\beta = -1$, all log masses equally likely) or with lower log masses more likely ($\beta = -2$, comparable to what is observed).

Note that the +1’s in the exponents in Equation (9) occur because for cluster_slug we specify the priors on the log of mass and age, while the conventional definitions of $\beta$ and $\gamma$ are in terms of distributions of mass and age, rather than log of mass and age.
and where we take the age distribution above $\sim$3 Myr to be either flat in age ($\gamma = 0$), flat in log age ($\gamma = -1$), or intermediate between the two ($\gamma = -0.5$). Flat in age is what would be expected if clusters, once formed, never disrupt, while flat in log age is what would be expected if clusters disperse such that the survival probability is equal for each decade in time. We do not consider variations in the prior on $A_V$, as there seems to be little theoretical or observational motivation to do so. The effects of varying the prior on $A_V$ are likely degenerate with the effects of varying the prior on age.

In Figures 9 and 10, we show how the masses and ages we derive for our star clusters depend on our choice of prior. We see that the choice of prior has relatively little impact on the results for most cluster masses, typically moving the median by an amount significantly less than the 10th–90th percentile range. The prior that deviates most from the others is $\beta = -2$, $\gamma = -1$. With this choice the majority of clusters are still largely unchanged, but a substantial tail of clusters appears for which the deviation $\Delta \log M$ is substantially less than zero, indicating that the prior $\beta = -2$, $\gamma = -1$ leads us to a substantially smaller mass estimate than alternative choices. The effect of prior choice on the inferred ages is somewhat greater, but as with masses, the only case that shows a very significant variation is $\beta = -2$, $\gamma = -1$, which again produces a tail of clusters with $\Delta \log T$ substantially negative. In most cases this is simply an understandable shift in power between two peaks of a bimodal PDF. As in the example shown above in Figure 8, in some cases there are two islands of probability that both fit the observed photometry reasonably well. In these cases a change in priors will enhance one island and suppress the other. Not surprisingly, in these cases the median shifts considerably.

However, there are also a number of extreme outliers, where the amount by which the posterior PDF shifts when we change our priors is so great that the previously inferred median now lies outside the 10th–90th percentile range. Many of these points represent clusters with substantial photometric errors, clusters that are not well-fit by our model libraries (perhaps because they are composites of multiple populations of different ages), or both. It is not surprising that the derived results for these clusters should be very sensitive to the choice of prior. When the error bars on the observations are large, the likelihood ratio is close to flat, and thus the posterior probability distribution is little altered from the prior one. Something analogous occurs if no model cluster in our library is a good fit to the observations, and instead there are a broad range of models that are less accurate fits.

However, there is also a population of clusters without large photometric error bars, and that are well-fit by our model library, that nevertheless show very substantial changes in their posterior PDFs depending on our choice of prior. To understand what is happening in these cases, in Figure 11 we show the 2D posterior PDF for an example cluster whose posterior PDF changes substantially depending on the choice of prior, but that does not have unusually large photometric errors, and
that is well-fit by our models. As the plot shows, this cluster also has two islands of probability, one centered at a mass of $\sim 10^{4.5} M_\odot$ and an age of $\sim 500$ Myr, and a second centered at a narrow range of ages $\sim 5-10$ Myr and masses of $10^2-10^3 M_\odot$. Unlike the degeneracy between age and extinction shown in Figure 8, these two possibilities need not sit at substantially different $A_V$ values. Instead, the degeneracy takes a rather different form, which is more closely related to IMF sampling.

The older age possibility corresponds to typical colors and luminosities for clusters of that mass and age range. On the other hand, the young age case corresponds to clusters that are on the far tail of the luminosity and color distributions for that age and mass. Thus the young age case represents extremely unusual photometric properties for clusters so young and low mass to have, resulting from very improbable draws from the IMF. If all ages are considered equally likely, then the rare
young, low-mass cases are rejected as unlikely. However, the difference between $\gamma = -1$ and $\gamma = 0$, as indicated by red and blue colors in Figure 11, amounts to the difference between a prior that says that there should be $\sim 100$ times as many clusters with ages from 5 to 10 Myr as clusters with ages of 505–510 Myr, and a prior that says there should be roughly equal numbers of clusters in those two ranges. If our prior is the former rather than the latter, then our Bayesian estimate assigns comparable total probabilities to the two possible fits, leading to a very different posterior PDF.

It is an open question to what extent the dependence on the choice of prior could be reduced or removed by the availability of H$_\alpha$ data, which is particularly sensitive to the youngest ages and therefore good at discriminating between otherwise-degenerate models. Fouesneau et al. (2012), analyzing star clusters in M83, found that H$_\alpha$ (as measured in their case by HST’s F657N filter) was very helpful in breaking degeneracies between fits. However, their analysis did not make use of F275W and instead had F336W as its bluest filter. The extra UV coverage provided by LEGUS’s F275W band should provide at least some of the same sensitivity to very young ages that Fouesneau et al. obtained from their H$_\alpha$ data. Moreover, Fouesneau et al.’s sample was limited to clusters with masses above $\sim 10^3 M_\odot$, where stochastic effects are somewhat less important than in our sample. Since H$_\alpha$ is produced primarily by the most massive stars, it is particularly vulnerable to stochastic effects, which might reduce its ability to discriminate between models for our data set. In any case, a survey to obtain H$_\alpha$ data for a subset of LEGUS galaxies is underway (HST-GO-13773, PI: R. Chandar). As those data become available

Figure 10. Same as Figure 9, but showing a comparison of ages rather than masses derived using different priors.
we will re-run our analysis pipeline using them, which should provide an answer to this question. It may also be possible to break degeneracies and remove dependence on the choice of prior using other discriminators. For example, the amount of scatter in surface brightness within the aperture (Whitmore et al. 2011) and the number of red stars in the vicinity of the cluster (Kim et al. 2012) have both been suggested as age-indicators. At present it is not clear how to build these into a Bayesian framework such as the one we have developed, as this would require extension of the likelihood function to include this information.

3.4. Dependence on Choice of Tracks, IMF, Metallicity, Extinction Curve, and Nebular Emission

We next examine the extent to which our results depend on our choice of tracks, IMF, metallicity, and extinction curve. In Figure 12 we show a comparison between results derived using our fiducial model, pad_020_kroupa_MW, and results derived using models with different extinction curves, metallicities, IMFs, and stellar tracks, and omitting nebular emission. In all these comparisons we use our fiducial priors, $\beta = -2$, $\gamma = -0.5$, but the results are comparable for other priors provided that we use the same prior for each library. Examining the figure, it is clear that the choice of IMF and extinction curve make almost no difference to the final posterior PDF. This is to be expected. For the IMF, the Chabrier (2005) and Kroupa (2001) IMFs we have tried differ mostly at the brown dwarf end, which has little impact on the integrated light properties even for old stellar populations. They also differ in that one is truncated at 100 $M_\odot$ and the other at 120 $M_\odot$, but those stars appear to be rare enough that the difference made to the integrated light by including or omitting them is smaller than the level of variation induced by IMF sampling effects. As for

**Figure 11.** Triangle plot for an example cluster (ID 320) in NGC 628e. Panels are similar to those in Figure 8. In the panels along the central diagonal showing the 1D marginal PDFs, the black line corresponds to the fiducial case ($\beta = -2, \gamma = -0.5$), while the blue and red lines refer to the alternate priors $\beta = -2, \gamma = 0$ and $\beta = -2, \gamma = -1$, respectively. In the three lower left panels, blue and red colors show the 2D PDF on a logarithmic intensity map (as indicated in the colorbar) for $\beta = -2, \gamma = 0$ and $\beta = -2, \gamma = -1$, respectively. Black contours show the 2D PDF for the fiducial case, with contours placed at log probability density values starting at -1 and increasing by 0.5 per contour. The 1D and 2D PDFs are all normalized to have unit integral.

**Figure 12.** Comparison of posterior PDFs derived using our fiducial model, pad_020_kroupa_MW, and various alternatives: from top to bottom, pad_020_chabrier_MW (Chabrier 2005 IMF instead of Kroupa 2001 IMF), pad_020_kroupa_SB (starburst instead of Milky Way extinction curve), pad_020_kroupa_MW_noneb (same as the default library, but omitting nebular emission), gen_014_kroupa_MW (Geneva stellar tracks at $Z = 0.014$ vs. Padova tracks at $Z = 0.02$), pad_004_kroupa_MW ($Z = 0.004$ instead of $Z = 0.02$), and pad_008_kroupa_MW ($Z = 0.008$ instead of $Z = 0.02$). In each panel, the fiducial model is on the x-axis, and the comparison is on the y-axis. The left column shows masses, and the right shows ages. Points plotted are the 50th percentile estimates of each quantity, and error bars (plotted for only a subset of the points to reduce clutter) indicate the 10th–90th percentile range. Different colors and plot symbols correspond to different fields, as indicated in the legend.
$A_V$, recall that the $A_V$ values we infer are generally modest. If there is little extinction, then the shape of the extinction curve matters little. The choice of IMF also makes little difference, which is, again, not surprising.

Whether we include nebular emission or not only makes a difference at young ages and low masses, though the latter is a selection effect—our flux limit is such that our low-mass clusters are exclusively young. In this range, we find that models excluding nebular emission produce systematically higher masses. This result is easy to understand: if we ignore the light produced by the nebula, then a higher stellar mass is required to produce the observed light. The effect of ignoring nebular emission on ages is more subtle, and again points to the importance of priors. We find that models excluding the nebular light do not produce 50th percentile ages below $\approx 3$ Myr, while our fiducial models show no such exclusion. The result for the no-nebular case is driven by the fact that the colors of stars are essentially constant at ages $\lesssim 3$ Myr, so the likelihood function we compute by comparing to the observed colors is flat over this age range. Combined our prior that all ages $T$ below $10^5$ years are equally likely, the posterior distribution with respect to log $T$ must therefore peak at larger $T$. On the other hand, when including nebular emission colors do evolve, at least mildly, at younger ages, and thus the likelihood function is not flat and we differentiate ages below 3 Myr. As a result, the 50th percentile values populate the full range of ages $\lesssim 3$ Myr.

The choice of tracks has modest but non-negligible effects. Compared to our fiducial case, the Geneva tracks seem to favor somewhat lower masses and ages. However, it is important to notice that the changes are largest for those clusters that have very significant uncertainties, so that the one-to-one line still falls mostly within the 10th–90th percentile range. In effect, in those cases where the colors are ambiguous and the posterior PDF is multi-peaked, changing from one set of tracks to another can favor or disfavor one of the two peaks.

The largest effect is from metallicity. Compared to Solar metallicity, the lower metallicity models favor substantially lower ages and masses for intermediate mass and age clusters, and slightly higher masses and ages for the youngest and lowest mass clusters. The effect at young ages and low masses appears to be similar to that seen in the no-nebular versus nebular comparison. This results from a complex series of effects of metallicity on the color: lower metallicity increases the ionizing flux and raises the nebular temperature, thereby increasing continuum and hydrogen recombination line emission. On average it also weakens emission from metal lines, though some lines that are particularly temperature-sensitive (e.g., [O m] $\lambda 5007$) may also get brighter. The resulting shift in colors appears to favor older ages and higher masses for the youngest and lowest mass clusters. The effect at higher mass and age comes largely from the interaction of colors with IMF sampling effects. Choosing a low metallicity tends to drive the model colors to the blue, beyond the point where they agree well with the observations. To match the observed colors requires driving the colors back to the red, and one way of doing this is to have a lower mass cluster that under-samples the massive end of the IMF, thereby producing redder colors.

Beyond the individual model comparisons, perhaps the most striking result shown in Figure 12 is how little difference the choice of models makes. The only one of the variations that we have examined that appears to make a substantial difference is using $Z = 0.004$, and, at the youngest ages perhaps, including or not including nebular emission. The primary reason for this insensitivity is that widths of the posterior PDFs for age and mass estimates are sufficiently large that they swamp systematics such as the choice of extinction curve, IMF, evolutionary tracks, and, at least over a limited range, metallicity. Perhaps these choices would become important if we had additional observational constraints beyond the five photometric bands to which we currently have access, but at present they are not the limiting factors in our knowledge.

### 3.5. Comparison of Stochastic and Non-stochastic Models

Our final analysis step is a comparison of the results we obtain using cluster_slug stochastic models to those we get from the deterministic Yggdrasil code. For this comparison, we use our fiducial library, pad_020_kroupa_MW, and prior distribution, $\beta = -2, \gamma = -0.5$. We can investigate this question on two levels. We can first ask how well cluster_slug and Yggdrasil models match the observed photometry, and whether one provides better matches of the observations than the other. We can then ask how the results we derive for cluster properties, and the nominal errors on those results, compare for the two methods.

#### 3.5.1. How Well Do cluster_slug and Yggdrasil Reproduce the Observed Photometry?

To investigate the consistency or lack thereof between the cluster_slug and Yggdrasil models and the observed photometry, we must first develop a statistic to quantify it. For cluster_slug, we parameterize the quality of agreement between the observations and the library using $D_{\text{form}}^2$, the normalized photometric distance between an observed cluster and its 5th closest neighbor in the cluster_slug library. In this form, exactly as in Section 3.1. For Yggdrasil, the best fitting model is determined via a $\chi^2$ minimization procedure. From the value of $\chi^2$ determined for the best fit (and the number of degrees of freedom in the model), we can compute the $Q$ statistic $Q(\chi^2)$, which formally is the probability that the value of $\chi^2$ would exceed the measured one even if the model were correct. Thus for example $Q(\chi^2) = 0.3$ means that, even if the best fit model were true, there is a 30% chance that each observation of that cluster would yield photometry for which the $\chi^2$ value would exceed our measured value. Thus values of $Q(\chi^2)$ near unity correspond to good fits, while ones with $Q(\chi^2) \ll 1$ indicate either that the error bars have been underestimated or that the best fitting model in the Yggdrasil library is probably not a good fit to the data.

Figure 13 shows the distribution of $Q(\chi^2)$ versus $D_{\text{form}}^2$ for our photometric catalog. In this plot, good fits for cluster_slug lie on the left side of the diagram, while poor fits lie on the right. Similarly, good fits for Yggdrasil lie at the top of the plot, and poor fits toward the bottom. The plot shows, consistent with Figure 3, that data are generally well-reproduced by the cluster_slug library. The solid line in the histogram of $D_{\text{form}}^2$ models shows the expected distribution for a library with $\ell = 6$ models within an observational error circle and $M = 1$ dimensional data (see Section 3.1); it is simply the binned, differential version of the cumulative distribution shown in Figure 3. The cluster_slug distribution of errors is roughly consistent with this
distribution, confirming that most clusters have good cluster_slug fits.

The same is not true of the Yggdrasil models. For a good fit the \(Q(S)\) statistic should be distributed so that a fraction \(f\) of models have \(Q(S) < f\). In fact, Figure 13 shows that there is a significant excess models at \(Q < 1\). The solid line in the histogram of \(Q\) values shows the distribution we would expect for a good fit, and there are clearly more small values of \(Q\) than expected, and significantly fewer values close to unity. In particular, nearly 10% of models have \(Q < 10^{-2}\), while for a good fit such large discrepancy should occur 0.1% of the time. In contrast, \(<5\%\) of clusters have \(D_s^{\text{norm}} > 2\), and only \(1\%\) have \(D_s^{\text{norm}} > 3\).

One possible explanation for the excess of clusters with \(Q < 1\) is that our observed catalog includes some objects that are not truly simple stellar populations, for example because they are blends of two clusters of different ages. However, we can immediately reject this as the dominant explanation, because such clusters should also be fit poorly by cluster_slug models, placing them in the lower right corner of the main panel of Figure 13. While there are a few objects of this sort, there are also a very large number of clusters for which Yggdrasil returns a \(Q\) value of \(\sim 10^{-2}\) or less, but cluster_slug nevertheless finds at least 5 simulated clusters whose photometry matches the observed values within 1 or 2 times the photometric errors. Indeed, even for the subset of clusters for which Yggdrasil produces \(Q < 10^{-2}\), the mean value of \(D_s^{\text{norm}}\) is 1.8. This is despite the fact that Yggdrasil is only fitting the colors (since the mass and thus the absolute luminosity are left as free parameters), while cluster_slug is fitting both the colors and the absolute magnitude. In contrast, there is no corresponding population of clusters in the upper right part of the diagram, which would indicate that Yggdrasil finds a good fit but cluster_slug does not.

Many of the clusters for which Yggdrasil does not find good fits are those for which cluster_slug assigns relatively low masses; the mean of clusterslug’s 50th percentile mass estimate for the subset of clusters for which Yggdrasil finds \(Q < 10^{-2}\) is 1000 \(M_\odot\), smaller than the 50th percentile mass for the entire sample by a factor of several. This strongly suggests that much of the problem is driven by Yggdrasil’s assumption of a well-sampled IMF, which fails in the low mass clusters. For these cases, exploring the full range of stochastic color variation is crucial to finding a good match to the data. However, there also some relatively massive clusters for which cluster_slug finds a good match but Yggdrasil does not. These may be cases where, despite the cluster’s overall relatively large mass, its colors are still significantly influenced by a small number of stars undergoing short-lived phases of stellar evolution. For example, a small number of He core stars undergoing blue loops might dominate the UV light budget of an otherwise old and red stellar population. For these rare phases of stellar evolution, the continuous assumption used in Yggdrasil may fail even for relatively massive clusters.

3.5.2. Comparison of cluster_slug and Yggdrasil Best Fits and Errors

Having discussed how well the models libraries match the data, we now ask how well the predicted cluster properties and the errors on those properties agree between the two methods. Figure 14 shows a comparison of the best-fit Yggdrasil values to the 50th percentile cluster_slug values. This is not quite a fair comparison, since in some cases the 50th percentile cluster_slug mass does not actually lie near a probability maximum. Nonetheless, the obvious alternative, plotting the cluster_slug point at the probability maximum, is no better. Some imprecision is inevitable when comparing a full posterior PDF to a single best fit.

Comparing masses in Figure 14, we notice that the models generally agree fairly well at high masses, but at low masses the Yggdrasil models on average tend to produce lower masses compared to cluster_slug’s 50th percentile, making the scatter about the 1–1 line somewhat asymmetric. Despite this, the mean masses for the population as a whole (indicated by the yellow circles in Figure 14) are quite similar.

The origin of this effect is almost certainly IMF under-sampling. For \(\sim 500 M_\odot\) clusters, the most common outcome by number is that the cluster will under-sample the massive end of the IMF, and thus will be under-luminous compared to what would be expected for a fully sampled IMF. Because Yggdrasil assumes full sampling it assumes a mass to light ratio that is too low, and ends up with a mass estimate that is too low as well. In contrast, cluster_slug correctly accounts for the sampling effect and assigns a broad PDF whose median value lies at higher masses than Yggdrasil’s best fit. The effect begins to appear at masses below roughly 103.5 \(M_\odot\), consistent with earlier analyses (e.g., Elmegreen 2000; Cerviño & Luridiana 2004, 2006; da Silva et al. 2012; Fouesneau et al. 2014). However, when averaging over a large population of clusters, there will be a few that happen to be over-luminous rather than under-luminous for their mass, so that the mean is the same as for fully sampled IMF. Consequently, while there is an asymmetric bias in the ages of individual clusters, estimates for the mean of the entire population as much less affected.

There are also systematic offsets in Yggdrasil and cluster_slug’s age and \(A_V\) distributions. Here the
Figure 14. Comparison of cluster masses (top panel), ages (middle panel), and extinctions (bottom panel) computed with Yggdrasil and cluster_slug. In all panels, the central point is plotted at the 50th percentile value returned by cluster_slug on the x axis, and the best fit returned by Yggdrasil on the y axis. Error bars, which we show only on a subset of points for clarity, indicate the 68% confidence interval for Yggdrasil, and the 16th–84th percentile range for cluster_slug. Point colors and shapes indicate the field for each cluster, and black dashed lines show the 1–1 relation. In the top two panels, large yellow points show the mean logarithmic mass and age, respectively, for both cluster_slug and Yggdrasil. Mean masses are computed by averaging the clusters in bins of 50th percentile cluster_slug mass from log(\(M_\odot\)) = 2–3, 3–4, 4–5, 5–6, and 6–7; ages are computed via an analogous procedure over bins from log(\(T_{\text{yr}}\)) = 6–7, 7–8, 8–9, and 9–10. In the middle panel, the alignments of the Yggdrasil fits at a set of common ages is an artifact of the fitting procedure.

4. PROPERTIES OF THE CLUSTER POPULATIONS

Having analyzed the properties of the individual clusters and their statistical and systematic uncertainties, we now seek to extend our analysis to the properties of the cluster population as a whole. Before commencing this exercise, we note that the results we derive in this section should not be taken as representative of the underlying properties of star clusters in our sample galaxies. Instead, we seek to derive the properties of those clusters that have made it into the LEGUS photometric catalog, which is the convolution of the true distribution of properties with the catalog selection. To convert these to intrinsic properties our analysis would have to be corrected for completeness, and such a correction is beyond the scope of the present work. In general we note that, because our catalog consists of visually inspected clusters, and the main criterion for being subject to visual inspection is exceeding a magnitude limit, our sample should be closer to magnitude-limited than mass-limited. This will tend to result in a steeper age distribution than is present in reality (e.g., Lamers 2009).

4.1. Methods

Our goal here is to derive the probability distribution functions for ages and masses of clusters, or their joint distribution, given the posterior PDFs for each individual cluster that we have derived in the previous section. Deriving this requires some care. The traditional method to derive the properties of the population, when each cluster has been fit using a traditional \(\chi^2\) approach, is to create bins in mass and
age, and then assign clusters to the bin corresponding to the best fit. The errors on each bin can then be computed from Poisson statistics. This is the method we will use to analyze the Yggdrasil results below. As usual, one must make a choice of how to place the bins, and we will consider two cases: one where the bins are placed uniformly in the log of mass or age, and one where they are placed so as to have a fixed number of clusters per bin. Note that the latter choice also maintains a fixed value of the relative error on the PDF; since the error for a bin containing $N$ clusters is $\sqrt{N}$, the ratio of the error to the mean is just $1/\sqrt{N}$, and thus a fixed value of $N$ implies a fixed fractional error.

This method is clearly not well-suited to our posterior PDFs, which are broad and multiply peaked. Moreover, we wish to avoid binning, since binning necessarily entails the loss of information. Instead, we proceed via a bootstrap resampling procedure. Consider the case of the mass distribution; distributions of age, extinction, or any combination of the three can be treated in an analogous fashion. Our central estimate for the mass distribution of the population as a whole is simply the sum of the posterior PDFs of mass for all the observed clusters. That is, in a galaxy containing $N$ clusters, for which we have determined a posterior probability distribution for mass $p_m$ for the $i$th cluster, our central estimate of the population PDF is simply

$$p_{\text{pop}}(m) = \frac{1}{N} \sum_{i=1}^{N} p_i(m).$$

(10)

To estimate the uncertainty on this estimate, we perform $N_\text{trials}$ trials. For each trial we draw $N$ clusters from our observed sample with replacement, meaning that we may draw the same cluster more than once. We then compute the population PDF $p_{\text{pop},i}(m)$ for that trial via Equation (10). Thus after $N_\text{trials}$ trials we have $N_\text{trials}$ population PDFs, and at each mass $m$ we have $N_\text{trials}$ estimates for the value of $p_{\text{pop},i}(m)$. We can then compute a 90% confidence interval on $p_{\text{pop}}(m)$ as simply the range from the 5th to the 95th percentile of these $N_\text{trials}$ estimates, and similarly for any other confidence interval of interest to us. The result of this procedure is therefore both a central estimate and a confidence interval on $p_{\text{pop}}(m)$ at each mass $m$, obtained without the need for binning, and using the full posterior PDFs from our Bayesian analysis of the individual clusters.

### 4.2. Marginal Mass and Age Distributions

Figure 15 shows the inferred marginal mass distributions for the entire population of star clusters, summing over all three of our fields. We emphasize that these plots show $p(M)$ rather than $p(m)$, so that a flat line corresponds to equal numbers of clusters per logarithmic mass interval. We see that, for our cluster_sim based modeling, the mass distribution is consistent with a power law $dn/dM \sim M^{-\beta}$, or perhaps very slightly steeper, over the mass range from roughly $10^{3.5} - 10^{5.5} M_\odot$. This result is essentially independent of the choice of model library or prior, and with relatively small statistical uncertainty based on bootstrap resampling. At lower masses the mass distribution flattens, almost certainly as a result of incompleteness, and at the lowest masses, $M \lesssim 10^3 M_\odot$, the inferred posterior distribution depends fairly strongly on the prior. This is a result of posterior PDFs for individual clusters in this mass range being quite broad, so that the total shape depends on the prior. In comparison, the choice of library matters little at low masses. Our central estimate for the shape of the mass function at the very high mass end,
approaching $10^6 M_{\odot}$, depends strongly on both the choice of library and the priors, but the statistical uncertainty is so large (due to the small number of clusters) that the bootstrap confidence intervals are mostly overlapping.

Figure 16 shows the corresponding marginal age distributions. As with Figure 15, we are plotting $p(\log T)$ rather than $p(T)$ (in contrast, for example, to the similar Figure 14 of Fouesneau et al. 2014), so that a flat line indicates a distribution where the cluster survival probability per decade in age is constant. Absence of cluster disruption would appear as a line of slope +1 on this plot. Based on the cluster_slug analysis, these are distributed roughly as $dN/dT \sim T^{-1}$ at ages from roughly $10^7$–$10^{8.5}$ years, with some dependence of the slope on the choice of prior and on the model library. The largest outlier is the result derived from the Geneva library, but this should probably be regarded as less reliable in the relevant age range due to its omission of TP-AGB stars.

The results with all models and libraries also show an excess of clusters at ages around $10^{6.5}$ years. This is partly driven by the choice of prior, which is flat in age $T$ (and thus rising in $\log T$) up to this age. We have argued that this is in fact the correct prior to use in this age range, since cluster disruption will be negligible for clusters this young. However, the bump at $10^{6.5}$ years might also be affected by completeness, which will preferentially suppress clusters at ages older than this due to the rapid decline in luminosity at older ages. At ages much below $10^{6.5}$ years, the results appear to be quite sensitive to the choice of metallicity and tracks, and the treatment of nebular emission. The feature driving this is almost certainly the variations in nebular emission, which is affected directly by metallicity, and indirectly (through the ionizing luminosity) by the choice of tracks. Eventually all models converge to our prior, $dN/dT \sim$ constant, because, at ages below 1 Myr, even if we include nebular emission then there is no change in cluster colors with age. Since the colors do not distinguish between ages <1 Myr, the posterior at such young ages simply reduces to the prior.

We also see a drop off in the cluster population at ages above $\sim10^{7.5}$ years. At least some of this drop must occur because these clusters’ low luminosities place them below our sensitivity limit, but there may also be a real decline in cluster numbers at older ages as well. If so, this would imply an increase in the rate of cluster disruption at ages around $10^{9.5}$ years. However, given our incompleteness, it would be premature to draw any conclusions at this point. The shape of this drop off depends somewhat sensitively on the choice of prior, with different but still physically plausible priors giving shapes in this age range that can be outside each others’ 90% confidence intervals. It does not depend strongly on the choice of library, again with the exception of the Geneva models, which produce quite a different result but are likely less reliable. On the other hand, it is possible that the Padova libraries may also have difficulties in this age range, since it is particularly susceptible to choices in how one treats the hard-to-model TP-AGB phase.

As noted above, to turn the distributions shown in Figures 15 and 16 into mass and age distributions for clusters in general will require completeness correction, a topic that we leave to a future paper. Moreover, the completeness correction will quite different for NGC 628 than for NGC 7793 due to their different distances ($\sim10$ Mpc versus $\sim3.5$ Mpc). Our goal is therefore not to derive the properties of the cluster populations in these two galaxies, but rather to understand how such determinations are influenced by the method used to convert the photometry into physical properties. We can already see from the results obtained thus far that the age distribution, at least in certain age ranges, will not be independent of the choice of prior, or the choice of library.

Comparing the cluster_slug results to the ones obtained using Yggdrasil, we find that the mass PDFs are roughly consistent within the error bars, though this depends to some extent on how one chooses to bin the Yggdrasil results, and on the prior that one chooses for the cluster_slug analysis. At small masses where the statistical uncertainties are largest, the Yggdrasil results appear to be closest to what one obtains using cluster_slug with a prior distribution $\beta = -2$, $\gamma = -1$. For the age distributions there is more discrepancy between the Yggdrasil and cluster_slug results. The Yggdrasil data show sharp drops in the number of clusters at around $10^{7.3}$ and $10^{8.1}$ years, with excesses near $10^{6.8}$ and $10^{7.9}$ years, though the prominence of these features depends on how one chooses to bin. The same structure in the Yggdrasil-determined ages is visible as the horizontal banding in Figure 14, and appears to result from the choice to assign a single best fit even in parts of color space where colors

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**Figure 16.** Same as Figure 15, but now showing the age distribution instead of the mass distribution. Black lines show different priors in age corresponding to $\gamma = 0, -0.5$, and $-1.0$ at ages above $10^{6.5}$ years.
provide relatively poor constraints on the true age. The cluster Slug models correctly capture the uncertainty in such regions by spreading out the posterior probability, and show no corresponding features. Instead, the population PDF is smooth. Nonetheless, the broad distribution of ages in the Yggdrasil models is not dissimilar from that obtained with cluster Slug. Apparently while the stochastic and deterministic methods produce relatively large differences on a cluster-by-cluster basis, they agree much more closely when analyzing the cluster population as a whole.

4.3. Joint Mass–Age Distribution

In addition to examining the mass and age distributions separately, we can examine them jointly. Figure 17 shows the joint distributions of mass and age summed over all three fields, marginalizing over the extinction. All the panels shown use our fiducial model, but we examine a range of priors. We also show the Yggdrasil results, binned in the $(\log M, \log T)$ plane.

Qualitatively, the cluster Slug results with all priors and the Yggdrasil results agree that the observed cluster population mostly lies along a band that extends from masses of $\sim10^4 M_\odot$ and ages of $\sim10^6$ years up to masses of $\sim10^8 M_\odot$ and ages of $\sim10^9$ years. This distribution is doubtless heavily influenced by the completeness of the survey, which would prevent us from detecting clusters that lie in the upper left of the diagram. On the other hand, the absence of clusters in the lower right corner, corresponding to massive, young systems, is a real feature of the distribution, as such clusters should be readily detectable.

Examining the results a bit more closely, we can notice that the choice of priors does influence the final distribution in non-negligible ways. As the priors on mass and age become flatter, $\beta = -1$ versus $\beta = -2$, and $\gamma = 0$ versus $\gamma = -1$, clusters tend to shift from the lower left end of the populated band, at low masses and young ages, upward to higher masses and ages. Thus priors that favor low mass and young age, $\beta = -2, \gamma = -1$, produce a prominent bump at masses of $10^{2.5} - 3.5 M_\odot$ and ages near $10^6$ years, while flatter priors, $\beta = -1, \gamma = 0$, remove this feature and distribute the clusters more evenly across a range of masses and ages. Flatter priors also produce a small island of probability at a mass of $\sim10^{6.5} M_\odot$ and an age of $\sim10^7$ years. This island is produced by a handful clusters that are not well fit by either the cluster Slug or Yggdrasil models, making the cluster Slug results for them quite sensitive to the prior.

The other very noticeable difference is between the Yggdrasil and cluster Slug models. Here the key feature is that the Yggdrasil-derived masses and ages are confined to a much smaller portion of the age-mass plane than the cluster Slug ones. This reflects the much smaller uncertainties that Yggdrasil assigns based on its $\chi^2$ fitting approach, as opposed to the Bayesian method of cluster Slug. This approach tends to concentrate the probability toward the peaks, suppressing the lower probability wings that are retained by using the full posterior.

5. SUMMARY AND CONCLUSIONS

We describe and investigate the performance of a Bayesian method to derive the masses, ages, and extinctions of star clusters observed using broadband photometry. Our sample data set for this analysis consists of 621 visually confirmed star clusters drawn from NGC 628 and NGC 7793. These galaxies are part of the LEGUS sample, which provides HST photometry in the NUV, U, B, V, and I bands. Our method, implemented based on the slug software suite and its cluster analysis tool cluster Slug, uses kernel density estimation coupled to implied conditional regression to return the full marginal posterior probability distributions of mass, age, and...

23 For example, the cluster for which Yggdrasil assigns the highest mass, which appears in the right-most pixel in the Yggdrasil panel of Figure 17, has a $Q$ parameter of $1.9 \times 10^{-4}$, indicating a very poor fit. The normalized photometric nearest-neighbor distance $D_{\text{NN}} = 3.0$ for our fiducial library, indicating that no cluster Slug models land closer than 3σ to the cluster's photometric properties; alternate libraries do no better. Given the strong disagreement with all sets of models, it seems likely that this object is not well described as a simple stellar population, or that all the model tracks we have available are deficient.

24 http://www.slugsps.com
extinction. This technique proves to be particularly useful in two regimes. One is for low-mass clusters, below $\sim 10^{3.5} \, M_\odot$, where the IMF is not fully sampled and the relationship between physical and photometric properties is therefore non-deterministic, leading to a broad posterior PDF that is not well characterized by a single best fit. The second case is for clusters that lie at locations in color space where there are significant degeneracies between mass, age, and extinction, so that there are multiple possible fits of comparable likelihood at distinct sets of physical parameters. In both of these regimes, a full posterior PDF provides a more reliable result, and a more realistic estimate of the errors, than methods that return a single best fit with a local error distribution around it. Indeed, by comparing our results to those produced by the non-stochastic cluster fitting code Yggdrasil, we find that the stochastic method removes some of the artifacts found in the deterministic one, and that it returns substantially larger error distributions.

At the level of individual clusters, our analysis technique proves to be insensitive, within the plausible range of variations, to our choice of metallicity, extinction curve, stellar IMF, and evolutionary tracks. On the other hand, because in many cases the observed photometry is consistent with more than one combination of mass, age, and extinction, our choice of prior probabilities does affect the final result. This sensitivity occurs because, when more than one mass-age-extinction combination is consistent with the observed photometry, the relative weight we assign to different “islands” of probability depends on our priors. This means that, for an individual cluster, before drawing any firm conclusions about its age, or to a less extent its mass, one should be careful to ensure that the results are robust against variations in priors, whether explicit or implicit. It is possible that the addition of extra photometric data, particularly Hα, might break some of these degeneracies, but that remains to be seen.

At the level of a cluster population, we find that the mass distribution in the range $10^3 - 10^6 M_\odot$ is very robust against changes in either prior or assumed stellar models, and to whether we derive the masses using stochastic cluster_slug models or deterministic Yggdrasil ones. At lower masses, priors begin to matter because the large amount of stochastic variation induced by incomplete IMF sampling means that photometry no longer provides strong constraints on mass, leaving the posterior close to the prior. At the highest masses stochasticity is unimportant, but depending on the assumed metallicity and choice of tracks, the inferred mass can vary by up to $\sim 0.5$ dex. When the distribution is dominated by a small number of clusters, this translates directly into variations in the overall distribution. Age distributions are somewhat less robust that mass ones. They are at least mildly dependent on the choice of prior at all ages, and at ages $\lesssim 3$ Myr when nebular emission contributes significantly to the light, they are very sensitive to metallicity, choice of tracks, and the assumed efficiency with which ionizing photons are converted to nebular light within the observational apertures. They are also quite dependent on the choice of prior at young ages.

The method we develop in this paper is executed in an automated pipeline, which is efficient enough that we can derive marginal posterior PDFs for the high confidence cluster catalog of 621 clusters in well under an hour on a workestation-level machine. This pipeline is available at http://www.slugsps.com, and will form the basis of a full stochastic cluster catalog derived from all the star clusters in the LEGUS sample. These will be expanded to include Hα data from the Hα-LEGUS program (PI: R. Chandar) as they become available. The resulting data set will provide an unprecedented sample of cluster properties, with realistic uncertainty distributions, from across the star-forming part of the Hubble sequence.

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