Cyclotomic Yokonuma–Hecke algebras are cyclotomic quiver Hecke algebras

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Abstract

We prove that cyclotomic Yokonuma–Hecke algebras of type A are cyclotomic quiver Hecke algebras and we give an explicit isomorphism with its inverse, using a similar result of Brundan and Kleshchev on cyclotomic Hecke algebras. The quiver we use is given by disjoint copies of cyclic quivers. We relate this work to an isomorphism of Lusztig.

Introduction

Iwahori–Hecke algebras appeared first in the context of finite Chevalley groups, as centralizer algebras of the induced representation from the trivial representation on a Borel subgroup. Since then, both their structure and their representation theory have been intensively studied. In particular, they have been defined independently as deformations of the group algebra of finite Coxeter groups. Further, connections with many other objects and theories have been established (this includes the theory of quantum groups, knot theory, etc.). Many variations and generalisations of the “classical” Iwahori–Hecke algebras have yet been defined. Among these, the following ones will catch our interest in this paper: Ariki–Koike algebras, Yokonuma–Hecke algebras and finally cyclotomic quiver Hecke algebras.

In their seminal paper, Ariki and Koike [ArKo] introduced and studied generalisations of Iwahori–Hecke algebras of type A and B: the so called Ariki–Koike algebras. It turns out that these algebras can be seen as cyclotomic quotients of affine Hecke algebras of type A, and also as deformations of the group algebra of the complex reflection group $G(\ell,1,n)$. Such deformations, in the general case of complex reflection groups, have been defined by Broué, Malle and Rouquier [BMR]: in that sense, Ariki–Koike algebras are the Hecke algebras associated with $G(\ell,1,n)$.

One of the most important results on the representation theory of Ariki–Koike algebras is Ariki’s categorification theorem [Ar] (proving a conjecture of Lascoux, Leclerc and Thibon [LLT]). This result implies that the decomposition matrices of such algebras can be computed using the canonical bases for quantum groups in affine type A. Partially motivated by this work, Khovanov and Lauda [KhLau1,KhLau2] and Rouquier [Rou] have independently defined the same algebra, known as quiver Hecke algebra or KLR algebra, in order to categorify quantum groups. In fact, they have shown that we have the following algebra isomorphism:

$$U^-_{\Lambda}(g) \cong [\text{Proj}(H)] = \bigoplus_{\beta \in \mathbb{Q}^+} [\text{Proj}(H_{\beta})]$$

where $U^-_{\Lambda}(g)$ is the integral form of the negative half of the quantum group $U_{\Lambda}(g)$ associated with a symmetrisable Cartan datum, with $A = \mathbb{Z}[q, q^{-1}]$, the set $\mathbb{Q}^+$ is the positive root lattice associated with the Cartan datum, the algebra $H = \bigoplus_{\beta \in \mathbb{Q}^+} H_{\beta}$ is the quiver Hecke algebra corresponding to this Cartan datum and $[\text{Proj}(H)]$ is the Grothendieck group of the additive category of finitely generated graded projective H-modules. A cyclotomic version of this theorem was conjectured in [KhLau1], the cyclotomic categorification conjecture, which was later proved by Kang and Kashiwara [KanKa]. More specifically, for each dominant weight $\Lambda$ the algebra $H$ has a cyclotomic quotient $H^\Lambda$ which categorifies the corresponding highest weight module $V(\Lambda)$.

A big step towards understanding these cyclotomic quiver Hecke algebras was made by Brundan and Kleshchev [BrKL] and independently by Rouquier [Rou]. The first two authors proved that cyclotomic Hecke algebras of type A are particular cases of cyclotomic quiver Hecke algebras; a similar result in the affine case has also been proved by Rouquier. Brundan and Kleshchev also noticed that...
the cyclotomic Hecke algebra inherits the natural grading of the cyclotomic quiver Hecke algebra, whose grading allows in particular to study the graded representation theory of cyclotomic Hecke algebras (see for example [BrKl2]). Moreover, they established a connection between the cyclotomic categorification theorem in type A for quiver Hecke algebras and Ariki’s categorification theorem.

On the other hand, Yokonuma [Yo] defined the Yokonuma–Hecke algebras in the study of finite Chevalley groups: they arise once again as centralizer algebras of the induced representation from the trivial representation, but now on a maximal unipotent subgroup. Their natural presentation in type A has been transformed since (see [Ju1, Ju2, JuKa, ChPA]), and the one we use here is given in [ChPou]. Similarly to Ariki–Koike algebras, Yokonuma–Hecke algebras of type A can be viewed as deformations of the group algebra of $G(d,1,n)$. This deformation, unlike in the Ariki–Koike case, “respects” the wreath product structure $G(d,1,n) \cong (\mathbb{Z}/d\mathbb{Z}) \wr S_n$. The representation theory of Yokonuma–Hecke algebras has been first studied by Thiem [Th1, Th2, Th3], while a combinatorial approach to this representation theory in type A has been given in [ChPA, ChPA2]. In this latter paper [ChPA2], Chlouveraki and Poulain d’Andecy introduced and studied generalisations of these algebras: the affine Yokonuma–Hecke algebras and their cyclotomic quotients, which generalise affine Hecke algebras of type A and Ariki–Koike algebras respectively. The interest in Yokonuma–Hecke algebras has grown recently: in [CJKL], the authors defined a link invariant from Yokonuma–Hecke algebras which is stronger than the famous (such as the HOMFLYPT polynomial) obtained from classical Iwahori–Hecke algebras.

The first aim of this paper is to show that cyclotomic Yokonuma–Hecke algebras are particular cases of cyclotomic quiver Hecke algebras, generalising thus the results of Brundan and Kleshchev [BrKl]; our goal will be achieved in Section 4 with Theorem 4.1. In fact, every known result on the cyclotomic quiver Hecke algebra can be applied to the cyclotomic Yokonuma–Hecke algebra: this includes the cyclotomic categorification theorem and the existence of a graded representation theory. In order to prove the main result Theorem 4.1, our strategy is to define inverse algebra homomorphisms, by constructing the images of the defining generators of the corresponding algebras; we proceed as in [BrKl]. In particular, we give an argument to avoid doing the calculations of [BrKl] again, see for instance Remark 2.7. In Section 5, as in [BrKl] we consider the degenerate case: we will define the degenerate cyclotomic Yokonuma–Hecke algebras and show that they are cyclotomic quiver Hecke algebras as well (Theorem 5.15). Finally, in Section 6 we relate our results to an isomorphism obtained in [Lu, JaPA, PA]. To that end, we first prove a general result on (cyclotomic) quiver Hecke algebras, when the quiver is the disjoint union of full subquivers. Although similar situations have already been studied in the literature (see for instance [SVV, Theorem 3.15] or [RSVV, Lemma 5.33]), the result we obtain in our context seems to be new and of independent interest.

We give now a brief overview of this article. Given a base field $F$ and $d,n \in \mathbb{N}^*$, we first define in Section 1 the cyclotomic Yokonuma–Hecke algebra $\widehat{\mathcal{Y}}_{d,n}(q)$ where $q \in F \setminus \{0,1\}$ has order $e \in \mathbb{N}_{>2} \cup \{\infty\}$ in $F^*$ and $\Lambda$ is a finitely-supported $e$-tuple of non-negative integers. We also define the quiver Hecke algebra $H_{\alpha}(Q)$ in full generality, where $\alpha$ is a composition of $n$ indexed by a set $\mathcal{K}$ and $Q = (Q_{\alpha,\beta})_{\alpha,\beta \in \mathcal{K}}$ is a matrix satisfying some properties. Considering particular cases for the matrix $Q = (Q_{\alpha,\beta})$, we define the cyclotomic quiver Hecke algebra $H^\Lambda_{\alpha}(\Gamma)$ where $\Lambda$ is now a finitely-supported tuple indexed by $\mathcal{K}$ and $\Gamma$ is a loop-free quiver without any multiple edge; in particular, with the exception of Section 6 we consider the case where $\alpha$ is given by $d$ disjoint copies of the (cyclic) quiver $\Gamma_e$ with $e$ vertices used in [BrKl]. We begin Section 2 by considering in $\widehat{\mathcal{Y}}_{d,n}(q)$ a natural system $\{e(\alpha)\}_{|\alpha| \leq d,n}$ of pairwise orthogonal central idempotents. Then, we define the “quiver Hecke generators” of $\widehat{\mathcal{Y}}_{d,n}(q) := e(\alpha)\mathcal{Y}_{d,n}(q)$ and we check that they verify the defining relations of $H^\Lambda_{\alpha}(\Gamma)$. In Section 3 we define the “Yokonuma–Hecke generators” of $H^\Lambda_{\alpha}(\Gamma)$ and again check the corresponding defining relations. We conclude the proof of the main theorem in Section 4 by showing that we have defined inverse algebra homomorphisms. We justify in Section 5 that the isomorphism of Theorem 4.1 remains true for the degenerate cyclotomic Yokonuma–Hecke algebra $\widehat{\mathcal{Y}}_{d,n}(1)$ that we define in §5.1. We end the section with Corollary 5.17, which states that, under some conditions, the algebras $\widehat{\mathcal{Y}}_{d,n}(q)$ and $\widehat{\mathcal{Y}}_{d,n}(1)$ are isomorphic. Finally, we begin Section 6 by some quick calculations about the minimal length representatives of the cosets of a Young subgroup in the symmetric group on $n$ letters $S_n$. The main results of the section are given in Theorems 6.26 and 6.30, where we prove an isomorphism about (cyclotomic) “disjoint quiver” Hecke algebra. We end the paper with Theorem 6.35: we show that we recover the isomorphism $\widehat{\mathcal{Y}}_{d,n}(q) \cong \oplus_{\lambda \vdash d,n} \text{Mat}_{m_\lambda,d,n} \widehat{H}_\Lambda(q)$ of [Lu, JaPA, PA], where $m_\lambda := \frac{\lambda_1!}{\lambda_1! \cdots \lambda_d!}$ and the algebra $\widehat{H}_\Lambda(q)$ is a tensor product of cyclotomic Hecke algebras.
1 Setting

Let $d, n \in \mathbb{N}^*$ and let $F$ be a field which contains a primitive $d$th root of unity $\xi$; in particular, the characteristic of $F$ does not divide $d$. We consider an element $q \in F^\times$ and we define $e \in \mathbb{N}^* \cup \{\infty\}$ as the smallest integer such that $1+q+\cdots+q^{e-1} = 0$. We will sometimes use the quantum characteristic of $F$, given by:

$$\text{char}_q(F) := \begin{cases} e & \text{if } e < \infty, \\ 0 & \text{if } e = \infty, \end{cases}$$

in particular $\text{char}_1(F)$ is exactly the usual characteristic of $F$. Except in Section 5, the element $q$ will always be taken different from 1. We set $I := \mathbb{Z}/\text{char}_q(F)\mathbb{Z}, J := \mathbb{Z}/d\mathbb{Z} \simeq \{1, \ldots, d\}$; unless mentioned otherwise, we have $K := I \times J$.

If $K$ is a set, we will refer to a finitely-supported tuple of non-negative integers $\Lambda = (\Lambda_k)_{k \in K} \in \mathbb{N}(K)$ as a weight. We say that a finitely-supported tuple $\alpha = (\alpha_k)_{k \in K} \in \mathbb{N}(K)$ is a $K$-composition of $n$ and we write $\alpha \vdash n$ if $(\alpha_k \neq 0$ for finitely many $k \in K$ and) $\sum_{k \in K} \alpha_k = n$. If $\alpha \vdash n$, we denote by $\mathcal{K}^{\alpha}$ the subset of $\mathcal{K}^n$ formed by the elements $k = (k_1, \ldots, k_n) \in \mathcal{K}^n$ such that:

$$\forall k \in \mathcal{K}, \# \{ a \in \{1, \ldots, n\} : k_a = k \} = \alpha_k,$$

that is, $(k_1, \ldots, k_n) \in \mathcal{K}^\alpha$ if and only if for all $k \in \mathcal{K}$, there are exactly $\alpha_k$ integers $a \in \{1, \ldots, n\}$ such that $k_a = k$. The subsets $\mathcal{K}^\alpha$ are the orbits of $\mathcal{K}^n$ under the natural action of the symmetric group on $n$ letters $\Sigma_n$, in particular each $\mathcal{K}^\alpha$ is finite.

1.1 Cyclotomic Yokonuma–Hecke algebras

Let $\Lambda = (\Lambda_i)_{i \in I} \in \mathbb{N}(I)$ be a weight; we assume that its level $\ell(\Lambda) := \sum_{i \in I} \Lambda_i$ verifies $\ell(\Lambda) > 0$. The cyclotomic Yokonuma–Hecke algebra of type $A$, denoted by $\tilde{Y}_{d,n}(q)$, is the unitary associative $F$-algebra generated by the elements

$$g_1, \ldots, g_{n-1}, t_1, \ldots, t_n, X_1$$

subject to the following relations:

$$t_a^d = 1, \quad t_a b = t_b t_a, \quad t_a g_a = g_a t_a, \quad g_a^2 = q + (q-1)g_a e_a, \quad g_a t_a = t_a g_a \quad \forall a - b > 1,$$

$$g_a+1 \cdot g_a+1 = g_a+1 \cdot g_a+1$$

where $s_a$ is the transposition $(a, a+1) \in \Sigma_n$ and $e_a := \frac{1}{d} \sum_{j \in J} t_a^j t_a^{-j}$, together with the following relations:

$$X_1 g_1, X_1 g_1 = g_1 X_1 g_1, \quad X_1 g_a = g_a X_1, \quad \forall a > 1,$$

$$X_1 t_a = t_a X_1,$$

and finally the cyclotomic one:

$$\prod_{i \in I} (X_i - q^i) = 0.$$  

Note that the presentation comes from [ChPA2], excepting the normalisation in (1.5) which was used in [ChPou]. In particular, it comes from (1.11) that $X_1$ is invertible in $\tilde{Y}_{d,n}(q)$. When $d = 1$, we recover the cyclotomic Hecke algebra $\tilde{H}_{A}(q)$ of [BrKl]; it is the cyclotomic Yokonuma–Hecke algebra $\tilde{Y}_{1,n}(q)$. In particular, the element $e_a$ becomes 1. We write $g_a^H$ (respectively $X_1^H$) for the element $g_a$ (resp. $X_1$) when $d = 1$, that is, considered in $\tilde{H}_{A}(q)$. 

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Following [ChPA2], we define inductively $X_{a+1}$ for $a \in \{1, \ldots, n - 1\}$ by

$$qX_{a+1} := g_a X_ag_a$$

(note that the $q$ comes from our different normalisation in (1.5)). As for $X_1$, we introduce the notation $X_a^\Lambda$ to denote $X_a$ in the case $d = 1$. The family $\{t_1, \ldots, t_n, X_1, \ldots, X_n\}$ is commutative and we have the following equalities:

$$g_aX_b = X_bg_a \quad \forall b \neq a, a + 1,$$

$$g_aX_{a+1} = X_ag_a + (q - 1)X_{a+1}e_a,$$

$$X_{a+1}g_a = g_aX_a + (q - 1)X_{a+1}e_a.$$  

The proof of the following result is the same as in [ChPA2, Proposition 4.7], where we write $g_w := g_{a_1} \cdots g_{a_r}$ for a reduced expression $w = s_{a_1} \cdots s_{a_r} \in S_n$; by Matsumoto’s theorem (see for instance [GeP6, Theorem 1.2.2]) the value of $g_w$ does not depend on the choice of the reduced expression, since the generators $g_a$ satisfy the same braid relations as the $s_a \in S_n$.

**Proposition 1.1.** The algebra $\tilde{\mathcal{Y}}_{d,n}(q)$ is a finite-dimensional $F$-vector space and a generating family is given by the elements $g_w X_{a_1}^\Lambda \cdots X_{a_n}^\Lambda t_{1}^{\alpha_1} \cdots t_{n}^{\alpha_n}$ for $w \in S_n$, $u_a \in \{0, \ldots, \ell(\Lambda) - 1\}$ and $v_a \in J$.

**Remark 1.2.** The above family is even an $F$-basis of $\tilde{\mathcal{Y}}_{d,n}(q)$, see [ChPA2, Theorem 4.15].

### 1.2 Cyclotomic quiver Hecke algebras

In their landmark paper [KhLa01], starting from a quiver without loop and multiple edges Khovanov and Lauda have constructed the so-called “quiver Hecke algebra”. Independently, Rouquier [Rou] made a similar construction, where the underlying object is a matrix $Q = (Q_{a,k'})_{k,k' \in \mathcal{K}}$ which contains the case of quivers. Here, we will first give this definition of [Rou], which generality will be used in Section 6 only, and then specialise to the case of quivers.

**1.2.1 General definition**

Let $\mathcal{K}$ be a set, $A$ a commutative ring, $u$ and $v$ two indeterminates and $Q = (Q_{a,k'})_{k,k' \in \mathcal{K}}$ a matrix satisfying the following conditions:

- the polynomials $Q_{a,k'} \in A[u,v]$ verify $Q_{a,k'}(u,v) = Q_{a',k}(v,u)$ for all $k, k' \in \mathcal{K}$;
- we have $Q_{k,k} = 0$ for all $k \in \mathcal{K}$.

Let $\alpha \models n$. The quiver Hecke algebra $H_\alpha(Q)$ associated with $(Q_{a,k'})_{k,k' \in \mathcal{K}}$ at $\alpha$ is the unitary associative $A$-algebra with generating set

$$\{e(k)\}_{k \in \mathcal{K}^n} \cup \{y_1, \ldots, y_n\} \cup \{\psi_1, \ldots, \psi_{n-1}\}$$

and the following relations:

$$\sum_{k \in \mathcal{K}^n} e(k) = 1,$$

$$e(k)e(k') = \delta_{k,k'}e(k),$$

$$y_a e(k) = e(k) y_a,$$

$$\psi_n y_b = y_b \psi_n \quad \text{if} \quad b \neq a, a + 1,$$

$$\psi_n \psi_b = \psi_b \psi_n \quad \text{if} \quad |a - b| > 1,$$

$$\psi_0 y_{a+1} e(k) = \begin{cases} (g_a \psi_a + 1)e(k) & \text{if} \quad k_a = k_{a+1}, \\ y_a \psi_a e(k) & \text{if} \quad k_a \neq k_{a+1}. \end{cases}$$

$$y_{a+1} \psi_a e(k) = \begin{cases} (g_a y_a + 1)e(k) & \text{if} \quad k_a = k_{a+1}, \\ \psi_a y_a e(k) & \text{if} \quad k_a \neq k_{a+1}. \end{cases}$$

$$\psi^2 e(k) = Q_{k_{a+1}, k_{a+2}}(y_a, 1 + y_a) e(k),$$

$$\psi_{a+1} \psi_a \psi_{a+1} e(k) = \begin{cases} \psi_a \psi_a \psi_a e(k) + \frac{Q_{k_a, k_{a+1}}(y_a, q_{a+1}) - Q_{k_{a+2}, k_{a+1}}(y_{a+2}, q_{a+1})}{y_{a+2} - y_{a+1}} e(k) & \text{if} \quad k_a = k_{a+2}, \\ \psi_a \psi_{a+1} \psi_a e(k) & \text{otherwise}. \end{cases}$$
1.2.2 Case of quivers

With the exception of Section 6, throughout this paper our matrix $Q$ will always be associated with a loop-free quiver without multiple edges and with vertex set $\mathcal{K}$. If $\Gamma$ is such a quiver, following [Rou, §3.2.4] we associate the following matrix $(Q_{k,k'})_{k,k'\in \mathcal{K}}$:

\[
Q_{k,k'}(u,v) := \begin{cases} 
0 & \text{if } k = k', \\
1 & \text{if } k \not\sim k', \\
v - u & \text{if } k \to k', \\
u - v & \text{if } k \leftarrow k', \\
-(u-v)^2 & \text{if } k \equiv k'. 
\end{cases}
\tag{1.34}
\]

where:

- we write $k \not\sim k'$ when $k \neq k'$ and neither $(k,k')$ or $(k',k)$ is an edge of $\Gamma$;
- we write $k \to k'$ when $(k,k')$ is an edge of $\Gamma$ and $(k',k)$ is not.
• we write \( k \leftarrow k' \) when \((k', k)\) is an edge of \( \Gamma \) and \((k, k')\) is not;
• we write \( k \rightleftharpoons k' \) when both \((k, k')\) and \((k', k)\) are edges of \( \Gamma \).

Moreover, we define:
\[
H_n(\Gamma) := H_n(Q),
\]
and if \( K \) is finite we also set \( H_n(\Gamma) := H_n(Q) \). Note that, in the setting of (1.34), the defining relations (1.26) and (1.27) become in \( H_n(\Gamma) \):

\[
\begin{align*}
\psi_a^2 e(k) &= \begin{cases} 
0 & \text{if } k_a = k_{a+1}, \\
\varepsilon(k) & \text{if } k_a \neq k_{a+1},
\end{cases} \\
(y_{a+1} - y_a)e(k) & \text{if } k_a \rightarrow k_{a+1}, \\
(y_a - y_{a+1})e(k) & \text{if } k_a \leftarrow k_{a+1}, \\
(y_{a+1} - y_a)(y_a - y_{a+1})e(k) & \text{if } k_a \rightleftharpoons k_{a+1},
\end{align*}
\]

\[
\psi_{a+1} \psi_a e(k) = \begin{cases} 
(\psi_a \psi_{a+1} \psi_a - 1) e(k) & \text{if } k_{a+2} = k_a \rightarrow k_{a+1}, \\
(\psi_a \psi_{a+1} \psi_a + 1) e(k) & \text{if } k_{a+2} = k_a \leftarrow k_{a+1}, \\
(\psi_a \psi_{a+1} \psi_a + 2y_{a+1} - y_a - y_{a+2}) e(k) & \text{if } k_{a+2} = k_a \rightleftharpoons k_{a+1}, \\
\psi_a \psi_{a+1} \psi_a e(k) & \text{otherwise}.
\end{cases}
\]

We now give a remarkable fact about quiver Hecke algebras; its proof only requires a simple check of the different defining relations.

**Proposition 1.7.** Let \( \Gamma \) be a loop-free quiver without multiple edges with vertex set \( K \). The quiver Hecke algebra \( H_n(\Gamma) \) is \( \mathbb{Z} \)-graded through:

\[
\begin{align*}
\deg e(k) &= 0, \\
\deg y_a e(k) &= 2, \\
\deg \psi_a e(k) &= -c_{k_a, k_{a+1}},
\end{align*}
\]

where \( C = (c_{k, k'})_{k, k' \in K} \) is the Cartan matrix of \( \Gamma \), defined by:

\[
c_{k, k'} := \begin{cases} 
2 & \text{if } k = k', \\
0 & \text{if } k \neq k', \\
-1 & \text{if } k \rightarrow k' \text{ or } k \leftarrow k', \\
-2 & \text{if } k \rightleftharpoons k'.
\end{cases}
\]

We now define the quivers which will be particularly important to us; we recall that \( I = \mathbb{Z}/\text{char}(F)\mathbb{Z} \) and \( J = \mathbb{Z}/d\mathbb{Z} \). We denote by \( \Gamma_e \) the following quiver:

• the vertices are the elements of \( I \);
• for each \( i \in I \) there is a directed edge from \( i \) to \( i+1 \).

In particular, for \( i, i' \in I \):

• we have \( i \rightarrow i' \) if and only if \( i' = i+1 \) and \( i \neq i' + 1 \);
• we have \( i \leftarrow i' \) if and only if \( i = i' + 1 \) and \( i' \neq i + 1 \);
• we have \( i \rightleftharpoons i' \) if and only if \( i = i' + 1 \) and \( i' = i + 1 \) (thus this only happens when \( e = 2 \));
• we have \( i \leftrightarrow i' \) if and only if \( i \neq i', i' \pm 1 \).

We give some examples in Figure 1. We now define the quiver

\[
\Gamma := \prod_{\varepsilon \in \varepsilon} \Gamma_e
\]

given by \( d \) disjoint copies of \( \Gamma_e \). Hence, our quiver \( \Gamma \) is described in the following way:

• the vertices are the elements of \( K := K = I \times J \);
• for each \((i, j) \in K \) there is a directed edge from \((i, j)\) to \((i + 1, j)\).
In particular, there is an arrow between \((i, j)\) and \((i', j')\) in \(\Gamma\) if and only if there is an arrow between \(i\) and \(i'\) in \(\Gamma_i\) and \(j = j'\). Moreover, the set \(K\) is finite if and only if \(e\) is finite.

We consider the diagonal action of \(\mathfrak{S}_n\) on \(K^n \simeq I^n \times J^n\), that is, \(\sigma \cdot (i, j) := (\sigma \cdot i, \sigma \cdot j)\). We will need the following notation:

\[
I^\alpha := \{i \in I^n : \exists j \in J^n, (i, j) \in K^\alpha\},
\]

\[
J^\alpha := \{j \in J^n : \exists i \in I^n, (i, j) \in K^\alpha\}.
\]

The sets \(I^\alpha\) and \(J^\alpha\) are finite and stable under the action of \(\mathfrak{S}_n\); note that \(K^\alpha\) is included in \(I^\alpha \times J^\alpha\) (we don’t have the equality in general).

Let now \(\Lambda := (\Lambda_k)_{k \in K} \in \mathbb{N}^{(K)}\) be a weight. The cyclotomic quiver Hecke algebra \(H^\Lambda_\alpha(\Gamma)\) is given by the quotient of the quiver Hecke algebra \(H_\alpha(\Gamma)\) by the relations:

\[
y^\Lambda_1 e(k) = 0, \quad \forall k \in K^\alpha; \quad (1.39)
\]

Note that this is indeed a particular case of (1.33). Note that the grading described in Proposition 1.7 is compatible with this quotient.

**Remark 1.8.** The cyclotomic Khovanov–Lauda algebra of \([\text{BrKl}]\) is the quiver Hecke algebra \(H^\Lambda_\alpha(\Gamma_e)\), that is, the algebra \(H^\Lambda_\alpha(\Gamma)\) for \(d = 1\). We write \(e_\Gamma^H, \ y^H_i, \ \psi^H_i\) the generators of \(H^\Lambda_\alpha(\Gamma_e)\); the reason for this notation will appear in \S 2.1.

The proof of the following result is the same as in \([\text{BrKl}, \text{Lemma 2.1}]\).

**Lemma 1.9.** The elements \(y_a \in H^\Lambda_\alpha(\Gamma)\) are nilpotent for \(a \in \{1, \ldots, n\}\).

As a corollary, together with Theorem 1.5 we recover a particular case of Theorem 1.6: in particular the algebra \(H^\Lambda_\alpha(\Gamma)\) is a finite-dimensional \(F\)-vector space.

**Remark 1.10.** Let us assume \(e < \infty\); the algebra \(H^\Lambda_\alpha(\Gamma)\) is hence finite-dimensional. However, comparing with Remark 1.2, it does not seem that easy to get its \(F\)-dimension; an answer will be given with Theorem 4.1 and (4.2).

## 2 Quiver Hecke generators of \(\widehat{Y}^\Lambda_\alpha(q)\)

Let \(\Lambda = (\Lambda_k)_{k \in K} \in \mathbb{N}^{(K)}\) be a weight; we assume that \(\ell(\Lambda) = \sum_{k \in K} \Lambda_k\) verifies \(\ell(\Lambda) > 0\). Moreover, we suppose that for any \(i \in I\) and \(j, j' \in J\), we have:

\[
\Lambda_{i,j} = \Lambda_{i,j'} := \Lambda_i.
\]

In particular, we will write \(\Lambda\) as well for the weight \((\Lambda_i)_{i \in I}\).

In this section, our first task is to define some central idempotents \(e(\alpha) \in \widehat{Y}^\Lambda_{d,n}(q)\) with \(\sum_{\alpha} e(\alpha) = 1\) for \(\alpha \models K\ n\). We will then prove the following theorem.

**Theorem 2.1.** For any \(\alpha \models K\ n\), we can construct an explicit algebra homomorphism:

\[
\rho : H^\Lambda_\alpha(\Gamma) \to \widehat{Y}^\Lambda_\alpha(q),
\]

where \(\widehat{Y}^\Lambda_\alpha(q) := e(\alpha)\widehat{Y}^\Lambda_{d,n}(q)\).

Note that \(\widehat{Y}^\Lambda_\alpha(q)\) is a unitary algebra (if not reduced to \(\{0\}\)), with unit \(e(\alpha)\). To define this algebra homomorphism, it suffices to define the images of the generators (1.16) and check that they verify the defining relations of the cyclotomic quiver Hecke algebra: the same strategy was used by Brundan and Kleshchev in \([\text{BrKl}]\) for \(d = 1\).

For a generator \(g\) of \(H^\Lambda_\alpha(\Gamma)\), we will use as well the notation \(q\) for the corresponding element that we will define in \(\widehat{Y}^\Lambda_\alpha(q)\). There will be no possible confusion since we will work with elements of \(\widehat{Y}^\Lambda_\alpha(q)\).
2.1 Definition of the images of the generators

We define now our different “quiver Hecke generators”.

2.1.1 Image of \(e(i,j)\)

Let \(M\) be a finite-dimensional \(\hat{\mathcal{Y}}_{d,n}^A(q)\)-module. Each \(X_a\) acts on \(M\) as an endomorphism of the finite-dimensional \(F\)-vector space (see Proposition 1.1); in particular, by (1.11) the eigenvalues of \(X_1\) can be written \(q^i\) for \(i \in \mathcal{I}\). Hence, applying [CuWa, Lemma 5.2] we know that the eigenvalues of each \(X_a\) are of the form \(q^i\) for \(i \in \mathcal{I}\). Concerning the \(t_a\), by (1.2) (they are diagonalizable and) their eigenvalues are \(d\)th roots of unity.

As the elements of the family \(\{X_a, t_a\}_{1 \leq a \leq n}\) pairwise commute, using Cayley–Hamilton theorem we can write \(M\) as the direct sum of its weight spaces (simultaneous generalized eigenspaces)

\[
M(i,j) := \{v \in M : (X_a - q^i)^N v = (t_a - \xi^j)v = 0 \text{ for all } 1 \leq a \leq n\}
\]  

(2.1)

for \((i,j) \in \mathcal{I}^n \times \mathcal{I}^n\), where \(N \gg 0\) and \(\xi\) is the given primitive \(d\)th root of unity in \(F\) that we considered at the very beginning of Section 1. Observe that some \(M(i,j)\) may be reduced to zero; in fact, only a finite number of them are non-zero.

**Remark 2.2.** The element \(e_a\) acts on \(M(i,j)\) as \(0\) if \(j_a \neq j_{a+1}\) and as \(1\) if \(j_a = j_{a+1}\).

We can now consider the family of projections \(\{e(k)\}_{k \in K^n}\) associated with the decomposition \(M = \bigoplus_{k \in K^n} M(k)\): the element \(e(k)\) is the projection onto \(M(k)\) along \(\bigoplus_{k' \neq k} M(k')\), in particular \(e(k)^2 = 0\) if \(k \neq k'\) then \(e(k)e(k') = 0\). Moreover, only a finite number of \(e(k)\) are non-zero.

As the \(e(k)\) are polynomials in \(X_1, \ldots, X_n, t_1, \ldots, t_n\) in fact \(e(k)\) is the product of commuting projections onto the corresponding generalized eigenspaces of \(X_a\) and \(t_a\), they belong to \(\hat{\mathcal{Y}}_{d,n}^A(q)\).

**Remark 2.3.** The above polynomials do not depend on the finite-dimensional \(\hat{\mathcal{Y}}_{d,n}^A(q)\)-module \(M\).

We are now able to define our central idempotents. We set, for \(\alpha \vdash_K n\):

\[
e(\alpha) := \sum_{k \in K^n} e(k)
\]

(since \(K^n\) is a \(\mathcal{G}_n\)-orbit, the element \(e(\alpha)\) is indeed central). Though we will not use this fact, we can notice that according to [PA, Corollary 3.2] and [LyMa], the subalgebras \(\hat{\mathcal{Y}}_{d,n}^A(q) := e(\alpha)\hat{\mathcal{Y}}_{d,n}^A(q)\) which are not reduced to zero are the blocks of the Yokonuma–Hecke algebra \(\hat{\mathcal{Y}}_{d,n}^A(q)\); see also [CuWa, §6.3]. For \(d = 1\), we recover the element \(e_a\) of [BrKl, (1.3)].

We will sometimes need the following elements:

\[
e(\alpha)(i) := \sum_{j \in \mathcal{I}^n} e(\alpha)e(i,j) = \sum_{j \in \mathcal{I}^n} e(i,j),
\]

\[
e(\alpha)(j) := \sum_{i \in \mathcal{I}^n} e(\alpha)e(i,j) = \sum_{i \in \mathcal{I}^n} e(i,j).
\]

(2.2)

For \(d = 1\), we recover with \(e(\alpha)(i)\) the element \(e(i)\) of [BrKl, §4.1]; we denote it by \(e^H(i)\). Finally, note that:

\[
e(\alpha)(i) \cdot e(\alpha)(j) = e(\alpha)(j) \cdot e(\alpha)(i) = e(i,j).
\]

From now, unless mentioned otherwise, we always work in \(\hat{\mathcal{Y}}_{d,n}^A(q)\); every relation should be multiplied by \(e(\alpha)\) and we write \(e(i)\) (respectively \(e(j)\)) for \(e(\alpha)(i)\) (resp. \(e(\alpha)(j)\)).

We give now a few useful lemmas.

**Lemma 2.4.** If \(1 \leq a < n\) and \(j \in \mathcal{I}^n\) is such that \(j_a \neq j_{a+1}\) then we have:

\[
g_a^2 e(j) = q e(j),
\]

\[
g_a X_a e(j) = X_a g_a e(j),
\]

\[
X_{a+1} g_a e(j) = g_a X_a e(j).
\]

**Proof.** We deduce it the from relations (1.5), (1.14) and (1.15) and from \(e_a e(j) = 0\) (since \(j_a \neq j_{a+1}\), see Remark 2.2). □

For the next lemma, we should compare with [JaPA, (15)].
Lemma 2.5. For $1 \leq a < n$ and $j \in J^n$ we have $g_a e(j) = e(s_a \cdot j)g_a$.

Proof. Let $M := \widehat{Y}_n^\Lambda(q)$. Given the relation (1.4), we see that $g_a$ maps $M(j)$ to $M(s_a \cdot j)$. Fix now $j \in J^n$ and let $j' \in J^n$ and $v \in M(j')$. If $j' = j$ then we get:

$$g_a e(j)v = e(s_a \cdot j)g_a v \ (= g_a v),$$

and if $j' \neq j$, since $g_a v \in M(s_a \cdot j')$ we have:

$$g_a e(j)v = e(s_a \cdot j)g_a v \ (= 0).$$

Hence, $e(s_a \cdot j)g_a$ and $g_a e(j)$ coincide on each $M'(j')$ for $j' \in J^n$ thus coincide on $M = \oplus_{j'} M(j')$ and we conclude since $M = \widehat{Y}_n^\Lambda(q)$ is a unitary algebra.

Corollary 2.6. Let $1 \leq a < n$ and $j \in J^n$; if $j_a = j_{a+1}$ then $g_a$ and $e(j)$ commute.

Remark 2.7 (About Brundan and Kleshchev’s proof - 1). Let $1 \leq a < n$; if $j \in J^n$ verifies $j_a = j_{a+1}$, the following relations are verified between $g_a e(j)$ and the $X_b e(j)$ for $1 \leq b \leq n$ in the unitary algebra $e(j)\widehat{Y}_n^\Lambda(q)e(j)$:

$$g_a^2 = q + (q - 1)g_a,$$

$$qX_{a+1} = g_a X_ag_a,$$

$$g_a X_b = X_ag_a \quad \forall a, b \neq a, a + 1,$$

$$g_a X_{a+1} = X_ag_a + (q - 1)X_{a+1},$$

$$X_{a+1}g_a = g_a X_a + (q - 1)X_{a+1}.$$
Remark \textbf{2.10} We first define some invertible elements $Q_a(i, j) \in F[[y_a, y_{a+1}]]^\times$ for $1 \leq a < n$ and $(i, j) \in K^\times$ by:

$$Q_a(i, j) := \begin{cases} 1 - q + q y_{a+1} - y_a & \text{if } i_a = i_{a+1}, \\ (y_a(i) - q y_{a+1}(i))/y_a(i) - y_{a+1}(i) & \text{if } i_a \not< i_{a+1}, \\ (y_a(i) - q y_{a+1}(i))/y_a(i) - y_{a+1}(i))^2 & \text{if } i_a \not> i_{a+1}, \\ q^{a+1}/(y_a(i) - y_{a+1}(i)) & \text{else}, \end{cases}$$

where $f_{a,j} \in \{1, q\}$ is given for $j_a \not= j_{a+1}$ by:

$$f_{a,j} := \begin{cases} q & \text{if } j_a < j_{a+1}, \\ 1 & \text{if } j_a > j_{a+1}, \end{cases}$$

with $<$ being a total ordering on $J = \mathbb{Z}/d\mathbb{Z} \cong \{1, \ldots, d\}$. For $d = 1$ or for $j_a = j_{a+1}$, the power series $Q_a(i, j)$ coincides with the definition given at [BrKl, (4.36)].

**Remark 2.11.** The scalar $f_{a,j}$ is only an artefact: if $q$ admits a square root $q^{1/2}$ in $F$, we can simply set $f_{a,j} := q^{1/2}$.

Finally, we give an easy lemma about these $f_{a,j}$.

**Lemma 2.12.** If $j_a \not= j_{a+1}$ then $f_{a,j} f_{a,s_a,j} = q$.

We introduce a notation for a power series $Q \in F[[y_1, \ldots, y_n]]$: if $\sigma \in S_n$ is a permutation, we denote by $Q^\sigma$ the power series $Q^\sigma(y_1, \ldots, y_n) := Q(y_{\sigma^{-1}(1)}, \ldots, y_{\sigma^{-1}(n)})$; note that $(Q^\sigma)^r = Q^{\sigma^r}$ (we get a right action).

We will use later the following properties verified by $Q_a(i, j)$ (see [BrKl, (4.35)]), where $Q^\sigma(i, j) := Q(i, j)^\sigma$.

**Lemma 2.13.** We have:

$$Q_a^{s_{a+1}}(i, j) = Q_a^{s_{a+1}}(s_a s_{a+1} \cdot (i, j)), \quad Q_a^{s_{a+1}}(i, j) = Q_a^{s_{a+1}}(s_a s_{a+1} \cdot (i, j)).$$

**Proof.** We check only the first equality, the second being straightforward considering $(i', j') := s_a s_{a+1} \cdot (i, j)$. For $i \in I$, let us write $y_a(i) := q'(1 - y_a)$; in particular, we have $y_a(i_a) = y_a(i)$. Noticing that $s_{a+1}s_a = (a, a + 2, a + 1)$, we get for $j_a = j_{a+1}$:

$$Q_a(s_{a+1}s_a \cdot (i, j)) = \begin{cases} 1 - q + q y_{a+1} - y_a & \text{if } i_a = i_{a+1}, \\ (y_a(i_a) - q y_{a+1}(i_a))/y_a(i_a) - y_{a+1}(i_a) & \text{if } i_a \not< i_{a+1}, \\ (y_a(i_a) - q y_{a+1}(i_a))/y_a(i_a) - y_{a+1}(i_a))^2 & \text{if } i_a \not> i_{a+1}, \\ q^{a+1}/(y_a(i_a) - y_{a+1}(i_a)) & \text{else}, \end{cases}$$

and we conclude using $y_a(i_a)^{s_{a+1}} = y_a(i_{a+1})$ and $y_{a+1}(i_{a+1})^{s_{a+1}} = y_{a+2}(i_{a+2})$. If $j_a \not= j_{a+1}$, we have $Q_a(s_{a+1}s_a \cdot (i, j)) = f_{a,s_a,j} = \begin{cases} q & \text{if } j_a < j_{a+1}, \\ 1 & \text{if } j_a > j_{a+1}, \end{cases}$ and this is exactly $f_{a+1,j}$.
We now introduce the following element of $\hat{\mathcal{Y}}^A(q)$:

$$
\Phi_a := g_a + (1 - q) \sum_{(i,j) \in \mathcal{K}^a, i_a \neq j_a = j_a+1} (1 - X_aX_{a+1}^{-1})^{-1} e(i,j) + \sum_{(i,j) \in \mathcal{K}^a, i_a = j_a+1} e(i,j),
$$

where $(1 - X_aX_{a+1}^{-1})^{-1} e(k)$ denotes the inverse of $(1 - X_aX_{a+1}^{-1})e(k)$ in $e(k)\hat{\mathcal{Y}}^A(q)e(k)$. Note that this element is indeed invertible, since for $k = (i,j)$ with $i_a \neq j_a+1$ its only eigenvalue $1 - q^{a-a+1}$ is non-zero, thanks to the definition of $I$. In particular, we have:

$$
\Phi_a e(j) = g_a e(j) \quad \text{if } j_a \neq j_a+1.
$$

For $d = 1$ we get the "intertwining element" defined in [BrKl, §4.2], and we write it $\Phi^H_a$.

**Remark 2.14.** Though we will not need this until Section 3, we define now the power series $P_a(i,j) \in F[[y_a, y_{a+1}]]$ for $1 \leq a < n$ and $(i,j) \in \mathcal{K}^a$ by:

$$
P_a(i,j) := \begin{cases} 
1 & \text{if } i_a = i_a+1, \\
(1 - q)(1 - y_a(i)y_{a+1}(i)^{-1})^{-1} & \text{if } i_a \neq i_a+1, \\
0 & \text{if } j_a \neq j_a+1.
\end{cases}
$$

For $d = 1$ or if $j_a = j_a+1$ we recover the definition given at [BrKl, (4.27)]. Moreover, we have the following equality:

$$
\Phi_a = \sum_{k \in \mathcal{K}^a} (g_a + P_a(k)) e(k). \tag{2.4}
$$

Indeed, the only non-obvious fact to check is $P_a(i,j)e(i,j) = (1 - q)(1 - X_aX_{a+1}^{-1})^{-1} e(i,j)$ if $i_a \neq i_a+1$ and $j_a = j_a+1$, but this is clear by (2.3). We will also use the following equality (the same one as in Lemma 2.13):

$$
P_{a+1}^*(i,j) = P_{a+1}^*(s_{a+1}s_a \cdot (i,j)). \tag{2.5}
$$

**Lemma 2.15.** We have the following properties:

$$
\begin{align*}
\Phi_a e(j) &= e(s_a \cdot j)\Phi_a, \tag{2.6} \\
\Phi_a e(i,j) &= e(s_a \cdot (i,j))\Phi_a, \tag{2.7} \\
\Phi_a X_b &= X_b\Phi_a, \quad \forall b \neq a, a+1, \tag{2.8} \\
\Phi_a y_b &= y_b\Phi_a, \quad \forall b \neq a, a+1, \tag{2.9} \\
\Phi_a Q_a(k) &= Q_a(k)\Phi_a, \quad \forall |b - a| > 1, \tag{2.10} \\
\Phi_a \Phi_b &= \Phi_b \Phi_a, \quad \forall |b - a| > 1. \tag{2.11}
\end{align*}
$$

**Proof.** We will use results from [BrKl, Lemma 4.1] (which is this lemma for $d = 1$).

(2.6) Using Lemma 2.5, it is clear if $j_a \neq j_a+1$ since then $\Phi_a e(j) = g_a e(j) = e(s_a \cdot j) g_a = e(s_a \cdot j)\Phi_a$.

(2.7) If $j_a = j_a+1$, we claim that the relation comes applying (2.6) and Remark 2.7 on the equality $\Phi_a^H e^H(i) = e^H(s_a \cdot i)\Phi_a^H$. If $j_a \neq j_a+1$ it follows directly from Lemma 2.8.

(2.8) Straightforward using (1.13) and since $e(i,j)$ are polynomials in $X_1, \ldots, X_n$.

(2.9) Using (2.2), (2.7) and (2.8) we get:

$$
\begin{align*}
\Phi_a y_b &= \sum_{i,j} (1 - q^{-ib} X_b) e(s_a \cdot (i,j))\Phi_a \\
&= \sum_{i} (1 - q^{-ib} X_b) e(s_a \cdot i)\Phi_a = y_b\Phi_a,
\end{align*}
$$

since $(s_a \cdot i)_b = i_b$.

(2.10) Since $Q_a(k) \in F[[y_b, y_{b+1}]]$ and $b \neq a, a+1$ and $b+1 \neq a, a+1$ it follows from (2.9).
Let us write \( \Phi'_a := \Phi_a - g_a \). Using (1.13) and Lemma 2.8 we get:

\[
\Phi'_a g_b = g_b \left( (1 - q) \sum_{i_{a+1} \neq j_{a+1}} (1 - X_a X_{a+1}^{-1})^{-1} e(s_b \cdot (i, j)) \right) + \sum_{i_{a+1} = j_{a+1}} e(s_b \cdot (i, j)) = g_b \Phi'_a,
\]

and exchanging \( a \) and \( b \) we get \( g_a \Phi'_b = \Phi'_b g_a \). Noticing that \( \Phi'_a \Phi'_b = \Phi'_b \Phi'_a \) (we don’t use here \(|b - a| > 1\)) and using (1.6), we get \( \Phi_a \Phi_b = (g_a + \Phi'_a)(g_b + \Phi'_b) = g_a g_b + \Phi'_a g_b + g_a \Phi'_b + \Phi'_a \Phi'_b = g_b g_a + g_a \Phi'_b + \Phi'_a g_b + \Phi'_a \Phi'_b = (g_b + \Phi'_b)(g_a + \Phi'_a) = \Phi_b \Phi_a \).

We are now ready to define our elements \( \psi_a \) for \( 1 \leq a < n \):

\[
\psi_a := \sum_{k \in K} \Phi_a Q_a(k)^{-1} e(k) \in \hat{Y}_a^\Lambda(q).
\]

As usual, we write \( \psi_a^\Lambda \) for \( \psi_a \) when \( d = 1 \), and this element \( \psi_a^\Lambda \) corresponds with the \( \psi_a \) of [BrKl, §4.3]. Note finally that for \( j \in J \) we have:

\[
\psi_a e(j) = f_{a,j}^{-1} g_a e(j) \quad \text{if} \quad j_a \neq j_{a+1}.
\]

### 2.2 Check of the defining relations

We now check the defining relations (1.17)–(1.25), (1.35)–(1.36) and (1.39) for the elements we have just defined. The idea is the following: when an element \( e(i, j) \) lies in a relation to check, if \( j_a = j_{a+1} \) then we get immediately the result by Remark 2.7 rewriting the same proof as [BrKl, Theorem 4.2], and if \( j_a \neq j_{a+1} \) then it will be easy (at least, easier than in [BrKl]) to prove the relation. Recall that we always work in \( \hat{Y}_a^\Lambda(q) \) (in particular every relation should be multiplied by \( e(\alpha) \)) and we write \( e(i) \) and \( e(j) \) without any \( (\alpha) \).

(1.39) We do exactly the same proof as in \( \hat{Y}_a^\Lambda(q) \). Let \((i, j) \in K^n\) and set \( M := \hat{Y}_a^\Lambda(q) \); recall that the action of \( X_1 \) on \( M \) is given by the action of \( e(\alpha) X_1 \). By (1.11) we have \( \prod_{i \in I} (X_1 - q^{\lambda_i}) = 0 \), hence:

\[
\prod_{i \in I} \left( X_1 - q^{\lambda_i} \right)^{1} e(i) = 0. \tag{2.12}
\]

As an endomorphism of \( M(i) \), the element \( (X_1 - q^{\lambda_i}) \) is invertible if \( i \neq i_1 \) since its only eigenvalue \( (q^{\lambda_i} - q^{\lambda_i}) \) is non-zero (note that \( \Lambda_i \) may be equal to 0). This means that there exist elements \( (X_1 - q^{\lambda_i}) \) such that \( (X_1 - q^{\lambda_i}) e(i) \cdot (X_1 - q^{\lambda_i})^{-1} e(i) = e(i) \). Hence, multiplying by all these inverses, the equation (2.12) becomes:

\[
(X_1 - q^{\lambda_i}) e(i) = 0.
\]

Finally, since \( y_i^{\lambda_i} = \sum_{i' \in I^n} (1 - q^{-\lambda_i} X_1) e(i') \) we obtain:

\[
y_i^{\lambda_i} e(i, j) = (1 - q^{-\lambda_i} X_1) e(i, j) = 0. \tag{1.17}
\]

(1.18) Idem.

(1.19) Straightforward since \( g_a \) and \( e(\alpha) \) both lie in the commutative subalgebra generated by \( i_1, \ldots, i_n \) and \( X_1, \ldots, X_n \).

(1.20) Straightforward by (2.7) and since \( Q_a(k') \) and \( e(k') \) commute with \( e(\alpha) \).
(1.21) True since \( \{X_a\}_a \) is commutative.

(1.22) True by (2.9).

(1.23) Let \( |a - b| > 1 \). We have, using (1.20), (2.10) and (2.11):

\[
\psi_a \psi_b = \sum_k \Phi_a Q_a(k)^{-1} e(k) \psi_b
\]

\[
= \sum_k \Phi_a Q_a(k)^{-1} \psi_b e(s_b \cdot k)
\]

\[
= \sum_k \Phi_a Q_a(k)^{-1} \Phi_b Q_b(s_b \cdot k)^{-1} e(s_b \cdot k)
\]

\[
= \sum_k \Phi_b Q_b(s_b \cdot k)^{-1} \Phi_a Q_a(k)^{-1} e(s_b \cdot k).
\]

Hence, noticing that (see Remark 2.10) \( Q_a(k) = Q_a(s_b \cdot k) \) and \( Q_b(s_b \cdot k) = Q_b(s_a s_b \cdot k) \) we get:

\[
\psi_a \psi_b = \sum_k \Phi_b Q_b(s_b \cdot k)^{-1} \psi_a e(s_b \cdot k)
\]

\[
= \sum_k \Phi_b Q_b(s_b \cdot k)^{-1} e(s_a s_b \cdot k) \psi_a
\]

\[
\psi_a \psi_b = \psi \psi_a.
\]

(1.24) First, if \( k = (i, j) \) verifies \( j_a = j_{a+1} \) then by Remark 2.7 we get from:

\[
\psi^H a y_{a+1} e^H(i) = \begin{cases} 
(y_a^H \psi^H_a + 1) e^H(i) & \text{if } i_a = j_{a+1}, \\
(y_a^H \psi^H_a) e^H(i) & \text{if } i_a \neq j_{a+1},
\end{cases}
\]

the following equality:

\[
\psi_a y_{a+1} e(i, j) = \begin{cases} 
(y_a \psi_a + 1) e(i, j) & \text{if } i_a = j_{a+1} \text{ and } j_a = j_{a+1}, \\
(y_a \psi_a) e(i, j) & \text{if } i_a \neq j_{a+1} \text{ and } j_a = j_{a+1}.
\end{cases}
\]

Hence it remains to deal with the case \( j_a \neq j_{a+1} \) (and no condition on \( i \)). Using Lemma 2.9 we get:

\[
\psi_a y_{a+1} e(i, j) = \Phi_a Q_a(i, j)^{-1} y_{a+1} e(i, j)
\]

\[
= f_{a,j} y_a y_{a+1} e(i, j)
\]

\[
= y_a \Phi_a Q_a(i, j)^{-1} e(i, j)
\]

\[
\psi_a y_{a+1} e(i, j) = y_a \psi_a e(i, j).
\]

Finally, we have proved:

\[
\psi_a y_{a+1} e(k) = \begin{cases} 
(y_a \psi_a + 1) e(k) & \text{if } k_a = k_{a+1}, \\
(y_a \psi_a) e(k) & \text{if } k_a \neq k_{a+1},
\end{cases}
\]

which is exactly (1.24).

(1.25) Similar.

Remark 2.16. Thanks to relations (1.22), (1.24) and (1.25), given \( f \in F[[y_1, \ldots, y_n]] \) and \( k \in K^n \) such that \( k_a \neq k_{a+1} \) we have \( f \psi_a e(k) = \psi_a f^a e(k) \). In particular, this holds if \( j_a \neq j_{a+1} \) with \( k = (i, j) \).
Once again, the result is straightforward if \( j_a = j_{a+1} \) using Remark 2.7. Let us then suppose \( j_a \neq j_{a+1} \); hence, necessarily we have \( k_a \neq k_{a+1} \) so we have to prove \( \psi_a^2 e(k) \).  We have:

\[
\begin{align*}
\psi_a^2 e(k) &= \psi_a e(s_a \cdot k) \\
&= \Phi_a Q_a(s_a \cdot k)^{-1} e(s_a \cdot k) \\
&= f_{a,sa,j}^{-1} g_a \Phi_a Q_a(k)^{-1} e(k) \\
\psi_a^2 e(k) &= (f_{a,sa,j} g_a) e(k).
\end{align*}
\]

Applying Lemmas 2.4 and 2.12 we find \( \psi_a^2 e(k) = e(k) \) (recall \( e(k) = e(i) e(j) = e(j) e(i) \)) thus we are done.

If \( j_a = j_{a+1} = j_{a+2} \), we get the result using Remark 2.7. Let us then suppose that we are not in that case: we have to prove \( \psi_{a+1}^2 \psi_{a+1} e(k) = \psi_{a+1}^2 \psi_{a} e(k) \). We will intensively use (1.20); note also that:

\[
\Phi_a e(i,j) = \begin{cases} 
(g_a + (1-g)(1-X_{aA_{a+1}}^{-1})^{-1}) e(i,j) & \text{if } i_a \neq i_{a+1} \text{ and } j_a = j_{a+1}, \\
(g_a + 1) e(i,j) & \text{if } i_a = i_{a+1} \text{ and } j_a = j_{a+1}, \\
g_a e(i,j) & \text{otherwise } (j_a \neq j_{a+1}),
\end{cases}
\]

and:

\[
\psi_a e(i,j) = \begin{cases} 
\Phi_a Q_a(i,j)^{-1} e(i,j) & \text{if } j_a = j_{a+1}, \\
f_{a,j}^{-1} g_a e(i,j) & \text{if } j_a \neq j_{a+1}.
\end{cases}
\]

It is convenient to introduce some notation. The couple \( (i,j) \) shall only be modified by the action of \( sa \) or \( sa_{a+1} \), hence we only write \( ((i_0,sa_{a+1},i_{a+2}),(j_a,j_{a+1},j_{a+2})) \) for \( (i,j) \). Moreover, for clarity we forget comas and only write the indexing, substituting 0 to a; thus, \( ((i_0,sa_{a+1},i_{a+2}),(j_a,j_{a+1},j_{a+2})) \) becomes \( ((012),(012)) \). Finally, as \( S_n \) acts diagonally on \( I \times J \) we can write \( (012) \) instead of \( ((012),(012)) \). Because an example beats lines of explanation, here is one: \( \psi_0 e(012) \) for \( \psi_0 e(s_a \cdot k) \).

**Case** \( j_0 = j_1 \neq j_2 \). Let us first compute \( \psi_0 \psi_0 \psi_0 e(012) \) and \( \psi_0 \psi_0 \psi_0 e(012) \). We have:

\[
\begin{align*}
\psi_0 \psi_0 \psi_0 e(012) &= \psi_0 e(012) \\
&= \psi_0 e(012) \psi_0 \\
&= \psi_0 e(012) \psi_0 e(012).
\end{align*}
\]

Since \( Q_1(201)^{-1} \in F[[y_1,y_2]] \) and recalling Remark 2.16 we get:

\[
\psi_0 \psi_0 \psi_0 e(012) = \Phi_1 e(012) \psi_0 Q_1^{-1} e(012) = \Phi_1 e(012) \psi_0 Q_0(012)^{-1} e(012).
\]

By Lemma 2.13 we have \( Q_1^{-1} e(012) = Q_1(201)^{-1} e(012) = (Q_0(012))^{-1} = Q_0(012) \). Hence:

\[
\psi_0 \psi_0 \psi_0 e(012) = \Phi_1 e(012) \psi_0 Q_0(012)^{-1} e(012).
\]

As:

\[
\psi_0 \psi_0 \psi_0 e(012) = \psi_0 \psi_1 \Phi_0 e(012) Q_0(012)^{-1},
\]

to have (1.36) it suffices to prove:

\[
\Phi_1 e(012) \psi_0 \psi_1 = \psi_0 \psi_1 \Phi_0 e(012).
\]

Let us distinguish two subcases.

- If \( i_0 \neq i_1 \) then:

\[
\Phi_1 e(012) \psi_0 \psi_1 = (g_1 + (1-g)(1-X_{aA_{a+1}}^{-1})^{-1}) \psi_0 e(012).
\]

Recalling (1.13) and Lemma 2.4 we get:

\[
\begin{align*}
\Phi_1 e(012) \psi_0 \psi_1 &= f_{0,0(021)}^{-1} (g_1 g_0 + (1-g) g_0 (1-X_{aA_{a+1}}^{-1})^{-1}) \psi_0 e(012) \\
&= f_{0,0(021)}^{-1} f_{0,1(021)}^{-1} (g_1 g_0 g_1 + (1-g) g_0 g_1 (1-X_{aA_{a+1}}^{-1})^{-1}) e(012).
\end{align*}
\]

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Using the braid relation (1.7) this becomes, recalling (2.7):
\[
\Phi_1 e(201) \psi_0 \psi_1 = f_{0,(021)}^{-1} f_{1,(012)}^{-1} \left( g_0 g_1 g_0 + (1-q) g_0 g_1 (1-X_0 X_1^{-1})^{-1} \right) e(012)
\]
\[
= f_{0,(021)}^{-1} f_{1,(012)}^{-1} g_0 g_1 (g_0 + (1-q) (1-X_0 X_1^{-1})^{-1}) e(012)
\]
\[
= f_{0,(021)}^{-1} f_{1,(012)}^{-1} g_0 g_1 \Phi_0 e(012)
\]
\[
\Phi_1 e(201) \psi_0 \psi_1 = f_{0,(021)}^{-1} f_{1,(012)}^{-1} g_0 g_1 (102) \Phi_0,
\]
and then, noticing that \( f_{1,(012)} = f_{1,(021)} \) and \( f_{0,(021)} = f_{0,(120)} \), we obtain:
\[
\Phi_1 e(201) \psi_0 \psi_1 = f_{0,(120)}^{-1} f_{1,(012)}^{-1} g_0 g_1 e(102) \Phi_0
\]
\[
= f_{0,(120)}^{-1} g_0 e(120) \psi_1 \Phi_0
\]
\[
\Phi_1 e(201) \psi_0 \psi_1 = \psi_0 \psi_1 \Phi_0 e(012),
\]
thus (2.13) is proved.

- If \( i_0 = i_1 \) then:
  \[
  \Phi_1 e(201) = (g_1 + 1) e(201),
  \]
  \[
  \Phi_0 e(012) = (g_0 + 1) e(012),
  \]
  thus with the same calculation as above (even easier) we get:
  \[
  \Phi_1 e(201) \psi_0 \psi_1 = f_{0,(021)}^{-1} f_{1,(012)}^{-1} (g_1 g_0 g_1 + g_0 g_1) e(012)
  \]
  \[
  = f_{0,(120)}^{-1} f_{1,(012)}^{-1} (g_0 g_1 g_0 + g_0 g_1) e(012) = \psi_0 \psi_1 \Phi_0 e(012),
  \]
so we got (2.13).

Until the end of the proof we use the same arguments as here, arguments which we will thus not recall.

**Case** \( j_0 \neq j_1 = j_2 \). Similar.

**Case** \( j_0 = j_2 \neq j_1 \). Once again we begin with the computation of \( \psi_1 \psi_0 \psi_1 e(012) \) and \( \psi_0 \psi_1 \psi_0 e(012) \).

We have:
\[
\psi_1 \psi_0 \psi_1 e(012) = \psi_1 \Phi_0 Q_0 (021)^{-1} e(021) \psi_1 = \psi_1 \Phi_0 e(021) \psi_1 Q_0 (021)^{-1},
\]
and:
\[
\psi_0 \psi_1 \psi_0 e(012) = \psi_0 \Phi_1 Q_1 (102)^{-1} e(102) \psi_0 = \psi_0 \Phi_1 e(102) \psi_0 Q_1 (102)^{-1}.
\]

Since \( Q_0 (021)^{-1} = Q_1 (102)^{-1} \), it suffices to prove:
\[
\psi_1 \Phi_0 e(021) \psi_1 = \psi_0 \Phi_1 e(102) \psi_0.
\]

Once again we distinguish two subcases.

- If \( i_0 \neq i_2 \) then:
  \[
  \psi_1 \Phi_0 e(021) \psi_1 = \psi_1 e(201) \Phi_0 \psi_1
  \]
  \[
  = f_{0,(021)}^{-1} g_1 \left( g_0 + (1-q)(1-X_0 X_1^{-1})^{-1} \right) e(012)
  \]
  \[
  = f_{0,(021)}^{-1} f_{1,(012)}^{-1} \left( g_1 g_0 g_1 + (1-q) g_0^2 (1-X_0 X_1^{-1})^{-1} \right) e(012)
  \]
  \[
  \psi_1 \Phi_0 e(021) \psi_1 = f_{0,(021)}^{-1} f_{1,(012)}^{-1} \left( g_1 g_0 g_1 + (1-q) g_0^2 (1-X_0 X_1^{-1})^{-1} \right) e(012).
  \]
Similarly, we find:
\[
\psi_0 \Phi_1 e(102) \psi_0 = \psi_0 e(120) \Phi_1 \psi_0
\]
\[
= f_{0,(120)}^{-1} g_0 \left( g_1 + (1-q)(1-X_1 X_2^{-1})^{-1} \right) \psi_0 e(012)
\]
\[
= f_{0,(120)}^{-1} f_{0,(012)}^{-1} \left( g_0 g_1 g_0 + (1-q) g_0^2 (1-X_0 X_2^{-1})^{-1} \right) e(012)
\]
\[
\psi_0 \Phi_1 e(102) \psi_0 = f_{0,(120)}^{-1} f_{0,(012)}^{-1} \left( g_0 g_1 g_0 + (1-q) g_0^2 (1-X_0 X_2^{-1})^{-1} \right) e(012),
\]
thus we conclude since \( f_{1,(012)} = f_{0,(012)} \) and \( f_{1,(021)} = f_{0,(120)} \) (we see it on this particular case or we can use Lemma 2.13).
- If \( i_0 = i_2 \) we get as above, with \( \alpha := f_{1,(201)}^{-1} f_{1,(012)}^{-1} f_{0,(120)}^{-1} f_{0,(012)}^{-1} \):

\[
\psi_1 \Phi_{0e}(021) \psi_1 = \alpha(g_1 g_9 g_1 + g_1^2) e(012)
= \alpha(g_1 g_9 g_1 + q) e(012) = \alpha(g_9 g_9 g_9 + q) e(012)
= \alpha(g_9 g_9 g_9 + g_9^2) e(012) = \psi_0 \Phi_{1e}(021) e(012).
\]

Case \( \#\{j_0, j_{a+1}, j_{a+2}\} = 3 \). We have \( j_a \neq j_{a+1} \) and \( j_a \neq j_{a+2} \) and \( j_{a+1} \neq j_{a+2} \) thus we get immediately:

\[
\psi_1 \psi_0 \psi_1 e(012) = f_{1,(201)}^{-1} f_{1,(012)}^{-1} f_{1,(102)}^{-1} f_{1,(012)}^{-1} f_{0,(120)}^{-1} f_{0,(021)}^{-1} g_9 g_9 g_9 e(012)
= f_{0,(120)}^{-1} f_{1,(102)}^{-1} f_{0,(120)}^{-1} f_{0,(021)}^{-1} g_9 g_9 g_9 e(012) = \psi_0 \psi_1 \psi_0 e(012),
\]

since \( f_{1,(201)} = f_{0,(012)}, f_{0,(021)} = f_{1,(102)} \) and \( f_{1,(102)} = f_{0,(120)} \).

3 Yokonuma–Hecke generators of \( H^\Lambda_\alpha(\Gamma) \)

Let \( \Lambda \) be a weight as in Section 2. The aim of this section is to prove the following theorem.

**Theorem 3.1.** For any \( \alpha \vdash K n \), we can construct an explicit algebra homomorphism:

\[
\sigma : \tilde{Y}^\Lambda_{d,n}(q) \to H^\Lambda_\alpha(\Gamma).
\]

Note that we do not consider yet \( \tilde{Y}^\Lambda_{d,n}(q) \). In particular, it suffices to define the images of the generators (1.1) and check if they verify the defining relations of the cyclotomic Yokonuma–Hecke algebra. As in Section 2, we use the same notation for a generator and its image.

3.1 Definition of the images of the generators

It is easier this time to define these images. First, since the elements \( y_1, \ldots, y_n \) are nilpotent (Lemma 1.9), we can consider power series in these variables. Hence, the quantities \( P_i(k), Q_{a}(k) \) and \( y_a(i) \) that we defined in §2.1 are also well-defined as elements of \( H^\Lambda_\alpha(\Gamma) \). We define finally as in (2.2) the elements \( e(i) \) and \( e(j) \) of \( H^\Lambda_\alpha(\Gamma) \) for \( i \in I^\alpha \) and \( j \in J^\alpha \).

We recall that \( \xi \) is a primitive \( d \)th root of unity in \( F \). Our “Yokonuma–Hecke generators” of \( H^\Lambda_\alpha(\Gamma) \) are given below.

\[
g_a := \sum_{k \in K^\alpha} (\psi_a Q_{\alpha}(k) - P_{\alpha}(k)) e(k) \quad \text{for } 1 \leq a \leq n
\]

\[
t_a := \sum_{j \in J^\alpha} \xi^{\alpha} e(j) \quad \text{for } 1 \leq a \leq n
\]

\[
X_a := \sum_{i \in I^\alpha} y_a(i) e(i) \quad \text{for } 1 \leq a \leq n
\]

As usual, we write \( g_a^H \) and \( X_a^H \) for the corresponding elements when \( d = 1 \): we recover the elements of [BrKl, §4.4].

**Remark 3.2 (About Brundan and Kleshchev’s proof - II).** This remark is similar to Remark 2.7. If \( j \in J^\alpha \) verifies \( j_a = j_{a+1} \) and if a relation in [BrKl, §4] involves only \( \psi_a^H, e^H(i) \) for \( i \in I^\alpha \) and \( y_b^H \) for \( 1 \leq b \leq n \), while its proof does not require any cyclotomic relation (1.39), then by the same proof, the same relation is satisfied between \( \psi_a e(j), e(i,j) \) and \( y_b e(j) \) in the unitary algebra \( e(j) H^\Lambda_\alpha(\Gamma) e(j) \).

If \( j \) verifies in addition \( j_{a+1} = j_{a+2} \), we will be able to add relations with \( \psi_{a+1}^H \), which we substitute by \( \psi_{a+1} e(j) \).

3.2 Check of the defining relations

As in §2.2, we will use Remark 3.2 when \( j_a = j_{a+1} \) to get the result from the same corresponding proof of [BrKl, Theorem 4.3], and when \( j_a \neq j_{a+1} \) we will need a few calculations.

\[ (1.2) \quad \text{Straightforward since } e(j^*) e(j^*) = \delta_{j,j^*} e(j) \text{ and } \xi^d = 1. \]
(1.3) Straightforward since \( e(j)e(j') = e(j')e(j) \).

(1.4) According to (1.17), it suffices to prove \( t_b g_a e(i, j) = g_a t_{s_a(b)} e(i, j) \) for every \((i, j) \in K^n\). For \((i, j) \in K^n\), we have, using (1.20):

\[
t_b g_a e(i, j) = t_b(\psi_a Q_a(i, j) - P_a(i, j)) e(i, j)
= t_b [e(s_a \cdot (i, j)) \psi_a Q_a(i, j) - e(i, j) P_a(i, j)]
= \xi^{(s_a)} g_a Q_a(i, j) e(i, j) - \xi^{(s_a)} g_a P_a(i, j) e(i, j),
\]

and:

\[
g_a t_{s_a(b)} e(i, j) = g_a \xi^{(s_a(b))} e(i, j) = \psi_a Q_a(i, j) \xi^{(s_a(b))} e(i, j) - P_a(i, j) \xi^{(s_a(b))} e(i, j).
\]

Thus, given our hypothesis on \((s_a \cdot j)b = j_{a+1}(b)\), we conclude since that expression is symmetric in \(s_a\) hence:

\[
P_a(i, j) \xi^{(s_a)} = P_a(i, j) \xi^{(s_a(b))},
\]

- It is clear if \(b \notin \{a, a+1\} \) since \(b = s_a(b)\).
- If \(b \in \{a, a+1\}\), it is clear if \(j_a = j_{a+1}\) and obvious if \(j_a \neq j_{a+1}\) since then \(P_a(i, j) = 0\).

(1.5) Let \((i, j) \in K^n\) and let us prove \(g_a^2 e(i, j) = (g + (q-1)g_a) e(i, j)\); summing over all \((i, j) \in K^n\) will conclude. If \(j_a = j_{a+1}\) then it is immediate applying Remark 3.2 on \((q^H)^2 = q + (q-1)q^H\) and left-multiplying by \(e(i)\), recalling \(e_a e(j) = e(j)\) and Corollary 2.6. If now \(j_a \neq j_{a+1}\), it is clear if \(a = 0\) which is trivial. If \(a > 0\) then:

\[
g_a^2 e(i, j) = g_a(g_a e(i, j) - P_a(i, j)) e(i, j)
= f_{a,j} g_a e(i, j)
= f_{a,j} g_a e(s_a \cdot (i, j))
= f_{a,j}(\psi_a Q_a(s_a \cdot (i, j)) - P_a(s_a \cdot (i, j))) e(i, j)
= g_a^2 e(i, j) = f_{a,j} f_{a,s_a} g_a e(i, j),
\]

hence we conclude using Lemma 2.12 and (1.35), since \(j_a \neq j_{a+1}\) implies \((i_a, j_a) /\sim (i_{a+1}, j_{a+1})\).

(1.6) Let us prove \(g_a g_b e(k) = g_b g_a e(k)\) for every \(k \in K^n\). By (1.22) the element \(\psi\) commutes with the elements \(P_a(k)\) and \(Q_a(k)\) of \(F[[y_a, y_{a+1}]]\). Moreover, \(Q_a(s_a \cdot k) = Q_a(k)\) and \(P_a(s_a \cdot k) = P_a(k)\), hence:

\[
g_a g_b e(k) = g_a(\psi_b Q_a(k) - P_a(k)) e(k)
= g_a(\psi_b Q_a(k) - \psi_a e(k) P_a(k))
= (\psi_b Q_a(k) - P_a(k)) \psi_a Q_a(k) e(k) - (\psi_a Q_a(k) - P_a(k)) P_a(k) e(k)
= (\psi_b Q_a(k)) Q_a(k) e(k) - \psi_a Q_a(k) P_a(k) e(k) - \psi_b Q_a(k) P_a(k) e(k) + P_a(k) P_a(k) e(k),
\]

and we conclude since that expression is symmetric in \(a\) and \(b\) (recalling (1.23)).

(1.7) Again it suffices to prove \(g_{a+1} g_a g_{a+1} e(i, j) = g_a g_{a+1} g_a e(i, j)\) for all \((i, j) \in K^n\). If \(j_a = j_{a+1} = j_{a+2}\) we get the result using Remark 3.2. Let us then suppose that we are not in that case. We will intensively use (1.20); recall the following fact:

\[
ge_a e(i, j) = \begin{cases} (\psi_a Q_a(i, j) - P_a(i, j)) e(i, j) & \text{if } j_a = j_{a+1}, \\ f_{a,j} \psi_a e(i, j) & \text{if } j_a \neq j_{a+1}. \end{cases}
\]

Finally, as during the proof of (1.36) in §2.2, we write for example \(g_0 e(012)\) instead of \(g_a e(s_a \cdot k)\).

Thus, given our hypothesis on \(j_0, j_1\) and \(j_2\) we have:

\[
\psi_1 \psi_1 \psi e(012) = \psi_0 \psi e(012).
\]
Case $j_0 = j_1 \neq j_2$. Let us first compute $g_1g_0g_1e(012)$ and $g_0g_1g_0e(012)$. We set $\alpha := f_{1, (012)}f_{0, (021)}$; we have:

$$
g_1g_0g_1e(012) = f_{1, (012)}g_1g_0e(021)\psi_1 = f_{1, (012)}f_{0, (021)}g_0(021)\psi_1 = \alpha\psi_1Q_1(021)\psi_0\psi_1e(012) - \alpha P_1(021)\psi_0\psi_1e(012)
$$

Similarly we have $P_1^{\alpha + 1}(021) = P_0(012)$ (see (2.5)). Hence we obtain, using (3.1) and noticing $f_{1, (012)} = f_{0, (120)}$ and $f_{0, (021)} = f_{1, (102)}$:

$$
g_1g_0g_1e(012) = \alpha\psi_0\psi_1\psi_0e(012)Q_0(012) = \alpha\psi_0\psi_1e(012)P_0(012)
$$

and let us prove:

$$
\psi_0^2e(012) = \psi_2^2e(012) = e(012).
$$

Hence, using (3.2), with $\alpha := f_{1, (012)}f_{1, (201)}$:

$$
g_1g_0g_1e(012) = f_{1, (012)}g_1(\psi_0Q_0(021) - P_0(021))e(021)\psi_1 = f_{1, (012)}g_1(\psi_0Q_0(021) - P_0(021))e(021)\psi_1
$$

and similarly we have

$$
\alpha\psi_1\psi_0e(012)Q_0^\alpha(021) = \alpha\psi_1e(012)Q_0^\alpha(021) - \alpha\psi_1e(012)P_0^\alpha(021)
$$

$$
= \alpha\psi_1\psi_0e(012)Q_0^\alpha(021) - \alpha\psi_1e(012)P_0^\alpha(021)
$$

Notice $f_{1, (012)} = f_{0, (02)}$ and $f_{1, (201)} = f_{0, (02)}$ we get finally:

$$
g_1g_0g_1e(012) = f_{0, (02)}g_0(\psi_1Q_1(102) - P_1(102))e(102)\psi_0 = g_{0, (02)}g_0(\psi_1Q_1(102) - P_1(102))e(102)\psi_0
$$

Note: Since for $a \in \{1, \ldots, n - 1\}$ it is clear that $X_{a+1} = X_aX_{a+1}$, it remains to prove that $qX_{a+1} = g_aX_a$, which we conclude taking $a = 1$. As we proved (1.5), it suffices to prove (1.4). Let $(i, j) \in K^n$ and let us prove:

$$
g_aX_{a+1}e(i, j) = \begin{cases} 
X_a g_a e(i, j) + (q - 1) X_{a+1} e(i, j) & \text{if } j_a = j_{a+1}, \\
X_a g_a e(i, j) & \text{if } j_a \neq j_{a+1}.
\end{cases}
$$

Again, we deduce the case $j_a = j_{a+1}$ from Remark 3.2. If $j_a \neq j_{a+1}$ we have, using (1.20) and (1.24):

$$
g_aX_{a+1}e(i, j) = q^{a+1}g_a e(i, j)(1 - y_{a+1}) + f_{a, j} e(i, j)
$$

$$
\equiv q^{a+1}f_{a, j} e(i, j) (1 - y_{a+1}) e(i, j) = q^{a+1}f_{a, j} e(i, j) e(s_a - (i, j)) e(i, j)
$$

$$
g_aX_{a+1}e(i, j) = X_a g_a e(i, j).
$$
We prove in fact (1.13), that is, \( g_a X_b = X_b g_a \) for \( b \neq a, a + 1 \). As \( y_b \) commutes with \( \psi_a \) by (1.22) we have, for any \( k \in K^n \) (where \( y_a(k) := y_a(i) \) with \( k = (i, j) \)):

\[
\begin{align*}
g_a X_b e(k) &= g_a e(k) y_b(k) \\
&= y_b(k) \psi_a Q_a(k) - P_a(k) e(k) \\
&= y_b(k) e(s_a \cdot k) \psi_a Q_a(k) - y_b(k) e(k) P_a(k) \\
g_a X_b e(k) &= X_b g_a e(k),
\end{align*}
\]

since \( y_b(k) e(s_a \cdot k) = q^{s_a t_b} - 1 = y_b e(s_a \cdot k) = X_b e(s_a \cdot k) \).

**Theorem 4.1.** We give now the main result of our paper; let \( H \) be the unitary algebra isomorphism with the cyclotomic Yononuma–Hecke algebra:

\[
\Lambda \overline{W} \text{e give now the main result of our paper; let } H \text{ be the unitary algebra isomorphism with the cyclotomic Yononuma–Hecke algebra.}
\]

\[
H_\Lambda(\Gamma) \iso \Lambda \overline{W}(q).
\]

As finitely many \( \Lambda \overline{W}(q) \) are non-zero, we deduce from this isomorphism that only finitely many \( H_\Lambda(\Gamma) \) are non-zero. Hence, as \( \Lambda \overline{W}_n(q) = \oplus_{\alpha | n} \Lambda \overline{W}(q) \), defining the cyclotomic quiver Hecke algebra of degree \( n \) by (recall (1.31)):

\[
H_\Lambda(\Gamma) := \bigoplus_{\alpha | n} H_\Lambda(\Gamma),
\]

we get the unitary algebra isomorphism with the cyclotomic Yononuma–Hecke algebra:

\[
H_\Lambda(\Gamma) \iso \Lambda \overline{W}_n(q).
\]

Recalling that the cyclotomic quiver Hecke algebra is naturally graded (Proposition 1.7), we get the following corollary.

**Corollary 4.2.** The cyclotomic Yononuma–Hecke inherits the grading of the cyclotomic quiver Hecke algebra.

Moreover, as we obtain a presentation of \( \Lambda \overline{W}_n(q) \) which does not depend on \( q \), we also get another one (see Corollary 5.16 for a slight improvement).

**Corollary 4.3.** Let \( q' \in F \setminus \{0, 1\} \). If \( \text{char}_q(F) = \text{char}_{q'}(F) \) then the algebras \( \Lambda \overline{W}_n(q) \) and \( \Lambda \overline{W}_n(q') \) are isomorphic.

Let us now prove Theorem 4.1. First, as we have a (non-unitary) algebra homomorphism \( \Lambda \overline{W}(q) \to \Lambda \overline{W}_n(q) \), by Theorem 3.1 we get an algebra homomorphism \( \Lambda \overline{W}(q) \to H_\Lambda(\Gamma) \), that we still call \( \sigma \). We will prove that \( \sigma : \Lambda \overline{W}(q) \to H_\Lambda(\Gamma) \) and \( \rho : H_\Lambda(\Gamma) \to \Lambda \overline{W}(q) \) (from Theorem 2.1) verify \( \sigma \circ \rho = \text{id}_{H_\Lambda(\Gamma)} \) and \( \rho \circ \sigma = \text{id}_{\Lambda \overline{W}(q)} \). Since these are algebra homomorphisms, it suffices to prove that they are identity on generators. To clarify the proof, let us add a \( \gamma \) on the quiver Hecke generators of \( \Lambda \overline{W}(q) \) and a \( \Gamma \) on the Yononuma–Hecke generators of \( H_\Lambda(\Gamma) \) (there isn’t any confusion possible with the former notation referring to the case \( d = 1 \) since we won’t use it any more).
Thus, we have $\sigma(\rho(y_a)) = y_a$ for all $1 \leq a \leq n$ and that $\sigma(\rho(y_a)) = y_a$ for all $1 \leq a \leq n$ and that $\sigma(\rho(y_a)) = y_a$ for all $1 \leq a \leq n$.

Let us start by finding the image of $e(k)$ by $\sigma \circ \rho$. By definition of $\rho$ we have $\rho(e(k)) = e^Y(k)$, so we have to prove $\sigma(e^Y(k)) = e(k)$. Let $M := H_{q}^N(\Gamma)$; the algebra homomorphism $\sigma$ gives $M$ a structure of $\mathbb{X}_{q}^N(q)$-module, finite-dimensional thanks to Theorem 1.6. If $M(k)$ denotes the weight space as in (2.1), by Remark 2.3 we know that the projection onto $M(k)$ along $\oplus_{k' \neq k} M(k')$ is given by $\sigma(e^Y(k))$. We prove that $e(k)$ is this projection too.

Let $(i, j) \in K^n$. For $1 \leq a \leq n$, we have $\sigma(X_a) = \sum_{i' \in I^n} (q^{i''} - q^{i''}y_a)e(i')$ so:

$$\sigma(X_a) - q^{i''} = \sum_{i' \in I^n} \left[(q^{i''} - q^{i''}y_a) - q^{i''}y_a\right] e(i').$$

Since $y_a$ is nilpotent, thanks to (1.17)–(1.18) we have:

$$\{ v \in M : (\sigma(X_a) - q^{i''}v = 0 \forall a \} = \sum_{i' \in I^n} e(i').$$

hence, for $N \gg 0$ we have:

$$M(i) := \{ v \in M : (\sigma(X_a) - q^{i''}v = 0 \forall a \} = e(i)M.$$ 

In a similar way we have $M(j) = e(j)M$ where $M(j) := \{ v \in M : (\sigma(X_a) - q^{i''}v = 0 \forall a \}$, thus:

$$M(k) = e(k)M.$$ 

Hence, as $\oplus_k M(k) = M$ we conclude that $e(k)$ is the desired projection and finally $e(k) = \sigma(e^Y(k))$.

The end of the proof is without any difficulty. We have:

$$\sigma(\rho(y_a)) = \sigma(y_a^Y) = \sum_{i \in I^n} [1 - q^{-i}a \sigma(X_a)\sigma(e^Y(i))$$

$$= \sum_{i \in I^n} [1 - q^{-i}a X_a^I] e(i)$$

$$= \sum_{i \in I^n} \left[1 - q^{-i}a \sum_{i' \in I^n} y_a(i') e(i)\right] e(i)$$

$$= \sum_{i \in I^n} [1 - q^{-i}a y_a(i)] e(i)$$

$$= \sum_{i \in I^n} [1 - q^{-i}a q^n (1 - y_a)] e(i)$$

$$\sigma(\rho(y_a)) = y_a.$$ 

Thus, we have $\sigma(Q_a^Y(k)) = Q_a(k)$ and $\sigma(P_a^Y(k)) = P_a(k)$, hence, recalling (2.4):

$$\sigma(\rho(\psi_a)) = \sigma(\psi_a^Y) = \sum_{k \in K^n} \sigma(\Phi_a)\sigma(Q_a^Y(k))\sigma(e^Y(k))$$

$$= \sum_{k \in K^n} \left( \sum_{k' \in K^n} \left[\sigma(g_a) + \sigma(P_a^Y(k'))\right] e(k')\right) Q_a(k)^{-1} e(k)$$

$$= \sum_{k \in K^n} \left( g_a^H + P_a(k)\right) Q_a(k)^{-1} e(k)$$

$$= \sum_{k \in K^n} \left( \left[ \sum_{k' \in K^n} (\psi_a Q_a(k') - P_a(k')) e(k')\right] + P_a(k)\right) Q_a(k)^{-1} e(k)$$

$$= \sum_{k \in K^n} \left[ (\psi_a Q_a(k) - P_a(k)) + P_a(k)Q_a(k)^{-1} e(k)\right]$$

$$\sigma(\rho(\psi_a)) = \psi_a.$$ 

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4.3 Proof of $\rho \circ \sigma = \text{id}_{\tilde{Y}_{d,n}^\Lambda(q)}$

This is even easier: we have to check $\rho(\sigma(g_a)) = g_a$ for $1 \leq a < n$ and $\rho(\sigma(X_n)) = X_n$ and $\rho(\sigma(t_a)) = t_a$ for $1 \leq a \leq n$. We have:

$$\rho(\sigma(g_a)) = \sum_{k \in K^n} [\psi_a^Y Q_a^Y(k) - P_a^Y(k)]e(k)$$

$$= \sum_{k \in K^n} [\Phi_a Q_a^Y(k)^{-1} Q_a^Y(k) - P_a^Y(k)]e(k)$$

$$= \sum_{k \in K^n} [\Phi_a - P_a^Y(k)]e(k)$$

$$\rho(\sigma(g_a)) = g_a.$$  

Recalling (2.3):

$$\rho(\sigma(X_a)) = \sum_{i \in I^n} y_a^Y(i)e^Y(i) = \sum_{i \in I^n} X_a e^Y(i) = X_a.$$  

Finally:

$$\rho(\sigma(t_a)) = \sum_{j \in J^n} \xi_a e^Y(j) = \sum_{j \in J^n} t_a e^Y(j) = t_a.$$  

The proof of Theorem 4.1 is now over.

5 Degenerate case

In this section, we extend the previous results to the case $q = 1$. In particular, we need to define a new “degenerate” cyclotomic Yokonuma–Hecke algebra. Many calculations are not written, since they are entirely similar to the non-degenerate case. Note the following thing: since the cyclotomic quiver Hecke algebra has no $q$ in its presentation, we do not need to define some new cyclotomic quiver Hecke algebra.

We recall that since $q = 1$ we have $e = \infty$. Let $\Lambda = (\Lambda_k)_{k \in K} \in \Pi(K)$ be a weight; we assume that $\ell(\Lambda) = \sum_{k \in K} \Lambda_k$ verifies $\ell(\Lambda) > 0$. Moreover, as in Section 2 we suppose that for any $i \in I$ and $j, j' \in J$, we have:

$$\Lambda_{i,j} = \Lambda_{i,j'} =: \Lambda_i.$$  

In particular, we will write $\Lambda$ as well for the weight $(\Lambda_i)_{i \in I}$.

5.1 Degenerate cyclotomic Yokonuma–Hecke algebras

We introduce here the degenerate cyclotomic Yokonuma–Hecke algebra; this algebra can be seen as the rational degeneration of the cyclotomic Yokonuma–Hecke algebra $\tilde{Y}_{d,n}^\Lambda(q)$.

The degenerate cyclotomic Yokonuma–Hecke algebra of type $A$, denoted by $\tilde{Y}_{d,n}^A(1)$, is the unitary associative $F$-algebra generated by the elements

$$f_1, \ldots, f_{n-1}, t_1, \ldots, t_n, x_1, \ldots, x_n$$

subject to the following relations:

$$t_a^d = 1,$$  

$$t_a t_b = t_b t_a,$$  

$$t_b f_a = f_a t_a(t_b),$$  

$$f_a^2 = 1$$  

$$f_a f_b = f_b f_a \quad \forall |a - b| > 1,$$  

$$f_a x_b = x_b f_a + e_a,$$  

$$f_a x_{a+1} = x_a f_a + e_a,$$  

$$f_a x_b = x_b f_a \quad \forall b \neq a, a + 1,$$  

$$x_a t_b = t_b x_a,$$

where $e_a := \frac{1}{d} \sum_{j \in J} t_b t_{a+1}^{-1}$, together with the following relations:

$$x_a x_b = x_b x_a,$$  

$$f_a x_{a+1} = x_a f_a + e_a,$$  

$$f_a x_b = x_b f_a \quad \forall b \neq a, a + 1,$$  

$$x_a t_b = t_b x_a.$$  

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and finally the cyclotomic one:

$$\prod_{i \in I} (x_1 - i)^{\lambda_i} = 0. \quad (5.12)$$

We obtained this presentation by setting $X_n = 1 + (q - 1)x_n$ in $\hat{\mathcal{Y}}_{d,a}(q)$, simplifying by $(1-q)$ as much as we can and then setting $q = 1$ ( according to the transformation made by Drinfeld [Dr] to define degenerate Hecke algebras). As in the non-degenerate case, the element $e_n$ verifies $e_n^a = e_n$ and commutes with $f_a$. Finally, note some consequences of (5.5) and (5.9):

$$x_{a+1} = f_ax_a + f_a e_a, \quad (5.13)$$
$$x_{a+1} f_a = f_ax_a + e_a. \quad (5.14)$$

When $d = 1$, we recover the degenerate cyclotomic Hecke algebra $\hat{H}_a(1)$ of [BrKl]; it is the degenerate cyclotomic Yokonuma–Hecke algebra $\hat{Y}_{d,a}(1)$. In particular, the element $e_n$ becomes 1, and $f_a$ (respectively $x_a$) is the element $s_a$ (resp. $x_b$) of [BrKl, §3].

We will use the following lemma (see [ChPA2, Lemma 2.15] for the non-degenerate case).

**Lemma 5.1.** For $u, v \in \mathbb{N}$ we have the following equalities:

$$f_ax_a x_{a+1} = x_ax_a f_a, \quad (5.15)$$
$$f_a x_{u+1} = x_{u+1} f_a + e_a \sum_{m=0}^{a-1} x_a \cdot x_{a+1}^{-m}, \quad (5.16)$$
$$f_a x_{a+1} f_a = x_{a+1} f_a - e_a \sum_{m=0}^{a-1} x_a \cdot x_{a+1}^{-m}, \quad (5.17)$$
$$f_a x_u x_{a+1} = \begin{cases} 
    x_u x_{a+1} f_a + e_a \sum_{m=0}^{u-1} x_a \cdot x_{a+1}^{-m} & \text{if } u \leq v, \\
    x_u x_{a+1} f_a - e_a \sum_{m=0}^{u-1} x_a \cdot x_{a+1}^{-m} & \text{if } u \geq v.
\end{cases} \quad (5.18)$$

**Proof.** We deduce (5.15) from different previous relations. The relations (5.16) and (5.17) can be proved by an easy induction. The equality (5.18) follows finally from these previous equalities. □

As the elements $g_a$ for $1 \leq a < n$ verify the same braid relations as the $s_a \in S_n$, for each $w \in S_n$ there is a well-defined element $g_w := g_{a_1} \cdots g_{a_r} \in \hat{Y}_{d,a}(1)$ which does not depend on the reduced expression $w = s_{a_1} \cdots s_{a_r}$.

**Proposition 5.2.** The algebra $\hat{Y}_{d,a}(1)$ is a finite-dimensional $F$-vector space and a family of generators is given by the elements $f_w x_1^{u_1} \cdots x_n^{u_n} t_1^{v_1} \cdots t_n^{v_n}$ for $w \in S_n$, $u_a \in \{0, \ldots, \ell(A) - 1\}$ and $v_a \in J$.

**Proof.** We use a similar method to [ArKo, OgPA]. As the unit element belongs to the above family, it suffices to prove that the $F$-vector space $V$ spanned by these elements is stable under (right-)multiplication by the generators of $\hat{Y}_{d,a}(1)$.

Let us consider $\alpha \colon f_w x_1^{u_1} \cdots x_n^{u_n} t_1^{v_1} \cdots t_n^{v_n}$ as in the proposition. By (5.2) and (5.3) the element $\alpha x_a$ remains in $V$. Moreover, writing (by (5.4) and (5.10)):

$$\alpha f_a = f_w x_1^{u_1} \cdots x_n^{u_n-1} (x_a x_{a+1}^{u_n} f_a) x_{a+2}^{u_n+2} \cdots x_n^{u_n} t_1^{v_1} \cdots t_n^{v_n},$$

and using (5.18) we conclude that $\alpha f_a \in V$, noticing that the element

$$x_1^{u_1} \cdots x_{a-1}^{u_{a-1}} (e_a x_a x_{a+1}^{u_n+1}) x_{a+2}^{u_n+2} \cdots x_n^{u_n} t_1^{v_1} \cdots t_n^{v_n}$$

belongs to $V$ for every $0 \leq u_a, u_{a+1} < \ell(A)$. Finally, according to (5.13), to prove that $\alpha x_a$ remains in $V$ it suffices now to prove that $\alpha x_1 \in V$, but this is clear by (5.8), (5.11) and (5.12). □

Let now $M$ be a finite-dimensional $\hat{Y}_{d,a}(1)$-module; it is a finite-dimensional $F$-vector space thanks to Proposition 5.2. By (5.12), the eigenvalues of $x_1$ on $M$ belong to $I$. We prove in Lemma 5.4 that all the $x_a$ have in fact their eigenvalues in $I$: this is the degenerate analogue of [CuWa, Lemma 5.2], which we used in §2.1.
Lemma 5.3. We have:

\[ x_a \phi_a = x_{a+1}, \]

\[ \phi_a^2 = (x_{a+1} - x_a - e_a)(x_a - x_{a+1} - e_a), \]

where \( \phi_a \) is the “intertwining operator” defined by:

\[ \phi_a := f_a(x_a - x_{a+1}) + e_a. \]

Proof. These are straightforward calculations. We have, using (5.9):

\[ x_a \phi_a = (f_{x_a+1} - e_a)(x_a - x_{a+1}) + x_a e_a \]

\[ = f_a(x_a - x_{a+1})x_{a+1} + x_{a+1} e_a \]

and:

\[ \phi_a^2 = f_a(x_a - x_{a+1})f_a(x_a - x_{a+1}) + 2f_a(x_a - x_{a+1})e_a + e_a \]

\[ = f_a(f_a x_{a+1} - e_a - f_a x_a - e_a)(x_a - x_{a+1}) + 2f_a(x_a - x_{a+1})e_a + e_a \]

\[ = (x_{a+1} - x_a)(x_a - x_{a+1}) + e_a \]

\[ \phi_a^2 = (x_{a+1} - x_a - e_a)(x_a - x_{a+1} - e_a). \]

Lemma 5.4. The eigenvalues of \( x_a \) belong to \( I \) for every \( 1 \leq a \leq n \).

Proof. We proceed by induction on \( a \). The proposition is true for \( a = 1 \); we suppose that it is true for some \( 1 \leq a < n \). Let \( \lambda \) be an eigenvalue of \( x_{a+1} \) (in a algebraic closure of \( F \)). As the family \( \{x_a, x_{a+1}, e_a\} \) is commutative, we can find a common eigenvector \( v \) in the eigenspace of \( x_{a+1} \) associated with \( \lambda \); we have \( x_a v = iv \) and \( e_a v = \delta v \) for some \( i \in I \) (by induction hypothesis) and \( \delta \in \{0,1\} \) (since \( e_a^2 = e_a \)). We distinguish now whether \( \phi_a v \) vanishes or not:

- if \( \phi_a v \neq 0 \), we get by Lemma 5.3:

\[ x_a(\phi_a v) = \phi_a(x_{a+1} v) = \lambda \phi_a v, \]

hence \( \lambda \) is an eigenvalue for \( x_a \) and by induction hypothesis we get \( \lambda \in I \);

- if \( \phi_a v = 0 \), by the same lemma we have:

\[ \phi_a^2 v = (\lambda - i - \delta)(i - \lambda - \delta)v = 0, \]

hence \( \lambda = i \pm \delta \in I \).

5.2 Quiver Hecke generators of \( \hat{Y}_{AI}^{A}(1) \)

We proceed as in Section 2: we define some central idempotents, then some “quiver Hecke generators” on which we check the defining relations of \( \hat{Y}_{AI}^{A}(1) \). The proofs are entirely similar to the non-degenerate case (even easier; note that once again the “hard work” has been made in [3rK1]), hence we won’t write them down. However, we will still define the different involved elements.

5.2.1 Image of \( e(i,j) \)

Let \( M \) be a finite-dimensional \( \hat{Y}_{AI}^{A}(1) \)-module. We know that the \( t_a \) are diagonalizable with eigenvalues in \( J \). Hence, recalling Lemma 5.4, we can write (recall that the family \( \{X_a, t_a\}_{1 \leq a \leq n} \) is commutative):

\[ M = \bigoplus_{(i,j) \in J^n \times J^n} M(i,j), \]

with:

\[ M(i,j) := \{ v \in M : (x_a - i_a)^N v = (t_a - t^N) v = 0 \text{ for all } 1 \leq a \leq n \}. \]
with $N \gg 0$; since $M$ is a finite-dimensional $F$-vector space, only finitely many $M(i, j)$ are non-zero. Considering once again the family of projections \{\(e(k)\)\}_{k \in K^n}$ associated with $M = \oplus_{k \in K^n} M(k)$, we define for $\alpha \in K^n$:
\[
e(\alpha) := \sum_{k \in K^n} e(k),
\]
and we set \(\hat{Y}_n(1) := e(\alpha)\hat{Y}_{d,n}(1)\). We can now define, for $i \in I^n$ and $j \in J^n$:
\[
e(\alpha)(i) := \sum_{j \in J^n} e(\alpha)(i, j),
\]
\[
e(\alpha)(j) := \sum_{i \in I^n} e(\alpha)(i, j).
\]
(5.19)
In particular, with $e(\alpha)(i)$ for $d = 1$ we recover the element $e(i)$ of [BrKi, §3.1].

From now on, unless mentioned otherwise we always work in \(\hat{Y}_n(1)\); every relation should be multiplied by $e(\alpha)$ and we write $e(i)$ and $e(j)$ without any $\alpha$.

**Lemma 5.5.** If $1 \leq a < n$ and $j \in J^n$ is such that $j_a \neq j_{a+1}$ then we have:
\[
f_a x_{a+1} e(j) = x_a f_a e(j),
\]
\[
x_{a+1} f_a e(j) = f_a x_a e(j).
\]

**Lemma 5.6.** For $1 \leq a < n$ and $j \in J^n$ we have $f_a e(j) = e(s_a \cdot j) f_a$. In particular, if $j_a = j_{a+1}$ then $f_a$ and $e(j)$ commute. Moreover, if $j_a \neq j_{a+1}$ then $f_a e(i, j) = e(s_a \cdot (i, j)) f_a$.

**Remark 5.7** (About Brundan and Kleshchev’s proof - III). Let $1 \leq a < n$; if $j \in J^n$ verifies $j_a = j_{a+1}$, when a proof in [BrKi, §3.3] needs only the elements $f_a$, $x_b$, $e(i)$ and the corresponding relations in $\hat{H}_n(1)$, we claim that the same proof holds in $e(j)\hat{Y}_n(1) e(j)$). We extend this claim to the case $j_a = j_{a+1} = j_{a+2}$.

### 5.2.2 Image of $y_a$

We define the following elements of $\hat{Y}_n(1)$ for $1 \leq a \leq n$:
\[
y_a := \sum_{i \in I^n} (x_a - i_a) e(i) \in \hat{Y}_n(1).
\]

When $d = 1$ we recover the elements defined in [BrKi, §3.3]. These elements are nilpotent: we will be able to make calculations in the ring $F[[y_1, \ldots, y_n]]$.

**Lemma 5.8.** For $j \in J^n$ such that $j_a \neq j_{a+1}$ we have:
\[
f_a y_{a+1} e(j) = y_a f_a e(j),
\]
\[
y_{a+1} f_a e(j) = f_a y_a e(j).
\]

### 5.2.3 Image of $\psi_a$

We first define some elements $p_a(i, j) \in F[[y_a, y_{a+1}]]$ for $1 \leq a < n$ and $(i, j) \in K^n$ by:
\[
p_a(i, j) := \begin{cases} 1 & \text{if } i_a = i_{a+1}, \\
(i_a - i_{a+1} + y_a - y_{a+1})^{-1} & \text{if } i_a \neq i_{a+1}, \\
0 & \text{if } j_a = j_{a+1},
\end{cases}
\]
and then some invertible elements $q_a(i, j) \in F[[y_a, y_{a+1}]]^\times$ for $1 \leq a < n$ and $(i, j) \in K^n$ by:
\[
q_a(i, j) := \begin{cases} 1 + y_{a+1} - y_a & \text{if } i_a = i_{a+1}, \\
1 - p_a(i, j) & \text{if } i_a \neq i_{a+1}, \\
(1 - p_a(i, j)^2)/(y_{a+1} - y_a) & \text{if } i_a \neq i_{a+1}, \text{ if } j_a = j_{a+1}, \\
1 & \text{if } i_a \neq i_{a+1}, \text{ if } j_a = j_{a+1}, \\
(1 - p_a(i, j))/(y_{a+1} - y_a) & \text{if } j_a \neq j_{a+1}.
\end{cases}
\]

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Remark 5.9. As in [BrK], the explicit expression of \( q_a(i, j) \) does not really matter; we only need some properties verified by these power series.

**Lemma 5.10.** We have:

\[
p_{a+1}^s(i, j) = p_a^{s+1}(s_{a+1} s_a \cdot (i, j)),
\]

\[
q_{a+1}^s(i, j) = q_a^{s+1}(s_{a+1} s_a \cdot (i, j)).
\]

We now introduce the following element of \( \hat{Y}_a(1) \):

\[
\varphi_a := f_a + \sum_{(i, j) \in K^a} (x_a - x_{a+1})^{-1} e(i, j) + \sum_{(i, j) \in K^a} e(i),
\]

where \((x_a - x_{a+1})^{-1} e(k)\) denotes the inverse of \((x_a - x_{a+1}) e(k)\) in \(e(k) \hat{Y}_a(1) e(k)\). In particular, we have:

\[
\varphi_a e(j) = f_a e(j) \quad \text{if } j_a \neq j_{a+1},
\]

\[
\varphi_a = \sum_{k \in K^a} (f_a + p_a(k)) e(k).
\]

Moreover, for \( d = 1 \) the element \( \varphi_a \) is the “intertwining element” defined in [BrK, §3.2].

**Lemma 5.11.** We have the following properties:

\[
\varphi_a e(j) = e(s_a \cdot j) \varphi_a,
\]

\[
\varphi_a e(i, j) = e(s_a \cdot (i, j)) \varphi_a,
\]

\[
\varphi_a x_b = x_b \varphi_a \quad \forall b \neq a, a + 1,
\]

\[
\varphi_a y_b = y_b \varphi_a \quad \forall b \neq a, a + 1,
\]

\[
\varphi_a q_b(k) = q_b(k) \varphi_a \quad \forall \mid b - a \mid > 1,
\]

\[
\varphi_a \psi_b = \psi_b \varphi_a \quad \forall \mid b - a \mid > 1.
\]

Our element \( \psi_a \) is defined for \( 1 \leq a < n \) by:

\[
\psi_a := \sum_{k \in K^a} \varphi_a q_a(k)^{-1} e(k) \in \hat{Y}_a(1).
\]

When \( d = 1 \) this element \( \psi_a \) corresponds to the \( \psi_a \) of [BrK, §3.3]. Note finally that for \( j \in J^a \) we have:

\[
\psi_a e(j) = f_a e(j) \quad \text{if } j_a \neq j_{a+1}.
\]

### 5.2.4 Check of the defining relations

**Theorem 5.12.** The elements \( y_1, \ldots, y_n, \psi_1, \ldots, \psi_{n-1} \) and \( e(k) \) for \( k \in K^a(\Gamma) \) verify the defining relations (1.39)–(1.36) of \( H^a(\Gamma) \).

The painstaking verification is exactly the same as in §2.2: we apply Remark 5.7 on the proof of [BrK], Theorem 3.2 for the cases \( j_a = j_{a+1} \), and when \( j_a \neq j_{a+1} \) then entirely similar (even the same) relations as in §2.2 are verified. Note two small differences with the proof in §2.2:

- we shall write \((x_a - x_b)\) instead of \((1 - q)(1 - X_a X_b^{-1})\);
- the elements \( f_{a, j} \) are equal to 1.

### 5.3 Degenerate Yokonuma–Hecke generators of \( H^a_0(\Gamma) \)

We proceed as in Section 3. Once again, the proofs are entirely similar to the non-degenerate case, hence we do not write them down.

First of all, since the elements \( y_1, \ldots, y_n \in H^a_0(\Gamma) \) are nilpotent we can consider power series in these variables. Hence, the quantities \( p_a(k), q_a(k) \) that we defined in §5.2.3 are also well-defined as elements of \( H^a_0(\Gamma) \). We define finally as in (5.19) the elements \( e(i) \) and \( e(j) \) of \( H^a_0(\Gamma) \) for \( i \in I^a \) and \( j \in J^a \).
We recall that $\xi$ is a primitive $d$th root of unity in $F$. Our “degenerate Yokonuma–Hecke generators” of $H^\Lambda_n(\Gamma)$ are given below.

$$f_a := \sum_{k \in K_a} (\psi_a q_a(k) - p_a(k)) e(k) \quad \text{for} \quad 1 \leq a < n$$

$$t_a := \sum_{j \in J^a} \xi^{a j} e(j) \quad \text{for} \quad 1 \leq a \leq n$$

$$x_a := \sum_{i \in I^a} (y_i + a_i) e(i) \quad \text{for} \quad 1 \leq a \leq n$$

When $d = 1$, the element $f_a$ (respectively $x_a$) is the element $s_a$ (resp. $x_a$) of [BrKl, §3.4].

Remark 5.13 (About Brundan and Kleshchev’s proof - IV). Let $1 \leq a < n$; if $j \in J^a$ verifies $j_a = j_{a+1}$, when a proof in [BrKl, §3.4] needs only the elements $\psi_a e(j), y_a e(j), e(i, j)$ and the corresponding relations in $H^\Lambda_n(\Gamma)$, we claim that the same proof holds in $e(j)H^\Lambda_n(\Gamma)e(j)$. We extend this claim to the case $j_a = j_{a+1} = j_{a+2}$.

Finally, similarly to §5.2.4 we have the following theorem. Once again the check of the various relations is exactly the same as in §3.2.

Theorem 5.14. The elements $f_1, \ldots, f_{n-1}, t_1, \ldots, t_n, x_1, \ldots, x_n$ verify the defining relations (1.2)–(1.11) of $\hat{Y}_{d,n}(1)$.

5.4 Isomorphism theorem

We give now the degenerate version of Theorem 4.1.

Theorem 5.15. There is a presentation of the degenerate cyclotomic Yokonuma–Hecke algebra $\hat{Y}_{d,n}^\Lambda(1)$ given by the generators (1.16) and the relations (1.17)–(1.25), (1.35)–(1.36) and (1.39), that is, we have an algebra isomorphism:

$$H^\Lambda_n(\Gamma) \cong \hat{Y}_{d,n}^\Lambda(1).$$

The proof of this theorem is entirely similar to the one of Theorem 4.1. In particular, by Theorem 5.12 we can define an algebra homomorphism $\rho : H^\Lambda_n(\Gamma) \to \hat{Y}_{d,n}^\Lambda(1)$ and by Theorem 5.14 we can define another algebra homomorphism $\sigma : \hat{Y}_{d,n}^\Lambda(1) \to H^\Lambda_n(\Gamma)$. From the inclusion $\hat{Y}_{d,n}^\Lambda(1) \subseteq \hat{Y}_{d,n}^\Lambda(1)$ we deduce an algebra homomorphism $\tilde{\sigma} : \hat{Y}_{d,n}^\Lambda(1) \to H^\Lambda_n(\Gamma)$. We prove then that $\rho$ and $\sigma$ are inverse homomorphisms, taking the images of the different defining generators.

Together with Theorem 4.1 we get the following corollaries (cf. [BrKl, Corollary 1.3]).

Corollary 5.16. If $q$ and $q'$ are two elements of $F^\times$ with $\text{char}_q(F) = \text{char}_{q'}(F)$ then $\hat{Y}_{d,n}^\Lambda(q)$ and $\hat{Y}_{d,n}^\Lambda(q')$ are isomorphic algebras.

Corollary 5.17. If $F$ has characteristic $\text{char}_q(F)$ then the cyclotomic Yokonuma–Hecke algebra $\hat{Y}_{d,n}^\Lambda(q)$ is isomorphic to its rational degeneration $\hat{Y}_{d,n}^\Lambda(q)$. This applies in particular when $F$ has characteristic 0 and $q$ is generic.

6 Another approach to the result

In [Lu, JaPA, PA], the authors have proved the following algebra isomorphism:

$$\hat{Y}_{d,n}^\Lambda(q) \cong \bigoplus_{\lambda \in \Lambda_m} \text{Mat}_{m} \hat{H}_{\lambda}^\Lambda(q), \quad (6.1)$$

with $F = C, e = \infty, \hat{H}_{\lambda}^\Lambda(q) := \hat{H}_{\lambda_1}^\Lambda(q) \otimes \cdots \otimes \hat{H}_{\lambda_d}^\Lambda(q)$, where we write $|_{\geq d}$ instead of $|_{\geq J}$ and where $m_\lambda$ are some integers (see (6.2)). We will see how we can relate this isomorphism to our previous work. To that extent, we prove in Theorem 6.26 a general result on quiver Hecke algebras in the case where the quiver is given by a disjoint union of full subquivers. The isomorphism is built from the following map (see (6.25)):

$$\Psi_{\lambda,\Psi} : w \mapsto (\psi_{\lambda,\Psi} w \psi_{\lambda,\Psi}^{-1}) E_{\lambda,\Psi},$$

for $w \in e(t_1)H_n(q)e(t_1)H_n(q)$, where $E_{\lambda,\Psi}$ are elementary matrices. In particular:

- in §6.1.2 we introduce the elements $\pi_i$, the minimal-length representatives of the right cosets of $S_n$ under the action of the Young subgroup $S_\lambda$, where $\lambda \models n$: this will lead to some calculations which will only be needed to explicit our homomorphism $\Psi_{\lambda,\Psi};$
- in §6.2 we will study the elements $\psi_{\lambda,n}$, and we will go on with the previous calculations.
6.1 Setting

Let \( K \) be a finite set; we recall that \( d, n \in \mathbb{N}^+ \) and \( J = \mathbb{Z}/d\mathbb{Z} \simeq \{1, \ldots, d\} \). We consider a partition of \( K \) into \( d \) parts \( K = \bigcup_{j \in J} K_j \). We recall that the left action of \( w \in \mathfrak{S}_n \) on tuples is given by \( w \cdot (x_1, \ldots, x_n) := (w x_{\sigma(1)}, \ldots, w x_{\sigma(n)}) \). We may use some elementary theory about Coxeter groups: we refer for instance to [GePf] or [Hum]. In particular, in that context we will denote by \( t \) the usual length function \( \mathfrak{S}_n \to \mathbb{N} \). Finally, let us mention that in this section, we will write \( t \) for the elements of \( J^n \).

6.1.1 Labellings and shapes

Let \( \lambda = (\lambda_j)_{1 \leq j \leq d} \vdash_d n \) be a \( d \)-composition of \( n \); recall that it means \( \lambda_j \geq 0 \) and \( \lambda_1 + \cdots + \lambda_d = n \). We define the integers \( \lambda_1, \ldots, \lambda_d \), given by \( \lambda_j := \lambda_1 + \cdots + \lambda_j \) for \( j \in J \). In particular, \( \lambda_1 = \lambda \) and \( \lambda_d = n \); we also set \( \lambda_0 := 0 \). From now on, the letter \( \lambda \) always stands for a \( d \)-composition of \( n \).

**Definition 6.1.** Let \( k \in \mathcal{K}^n \) and \( t \in J^n \).

- We say that \( k \) is a labelling of \( t \) when the following rule is satisfied:
  \[
  \forall a \in \{1, \ldots, n\}, k_a \in K_{k_a},
  \]
  that is:
  \[
  \forall a \in \{1, \ldots, n\}, \forall j \in J, k_a \in K_j \iff t_a = j.
  \]
  We write \( \mathcal{K}^t \) for the elements \( \mathcal{K}^n \) which are labellings of \( t \).

- We say that \( t \) has shape \( \lambda \vdash_d n \) and we write \([t] = \lambda\) if for all \( j \in J \) there are exactly \( \lambda_j \) components of \( t \) equal to \( j \), that is:
  \[
  \forall j \in J, \# \{ a \in \{1, \ldots, n\} : t_a = j \} = \lambda_j.
  \]
  We write \( J^\lambda \) for the elements \( J^n \) with shape \( \lambda \).

The sets \( J^\lambda \) are the exact orbits of \( J^n \) under the action of \( \mathfrak{S}_n \), in particular \([w \cdot t] = [t]\) for every \( w \in \mathfrak{S}_n \) and \( t \in J^n \). Moreover, the cardinality of \( J^\lambda \) is:

\[
m_\lambda := \frac{n!}{\lambda_1! \cdots \lambda_d!}.
\]

We write \( t^\lambda \in J^\lambda \) for the trivial element of shape \( \lambda \), given by:

\[
\forall a \in \{1, \ldots, n\}, \forall j \in J, t^\lambda_a = j \iff \lambda_{j-1} < a \leq \lambda_j,
\]

that is:

\[
t^\lambda := (1, \ldots, 1, d, \ldots, d),
\]

where each \( j \in J \) appears \( \lambda_j \) times. Note that \( \mathcal{K}^{t^\lambda} \simeq \mathcal{K}^1 \times \cdots \times \mathcal{K}^{\lambda_d} \).

6.1.2 Young subgroups

Most results of this section are well-known; however, since in the literature they are stated either for a left or a right action (see Remark 6.9), for the convenience of the reader we state all of them with a left action. We remind the reader that some calculations made here will only be used in §6.4.3, namely with Lemmas 6.12 and 6.13.

Let \( \lambda \vdash_d n \); the following group:

\[
\mathfrak{S}_\lambda := \mathfrak{S}_{\lambda_1} \times \cdots \times \mathfrak{S}_{\lambda_d},
\]

can be seen as a subgroup of \( \mathfrak{S}_n \) (the “Young subgroup”), where we consider that \( \mathfrak{S}_\lambda \simeq \mathfrak{S}(\{\lambda_j-1 + 1, \ldots, \lambda_j\}) \). Recall that:

- the group \( \mathfrak{S}_\lambda \) (resp. \( \mathfrak{S}_{\lambda_j} \)) is generated by \( s_1, \ldots, s_{\lambda_j-1} \) (resp. \( s_{\lambda_j-1+1}, \ldots, s_{\lambda_j} \))
- the subgroup \( \mathfrak{S}_\lambda \) is generated by all the \( s_a \) for \( a \in \{1, \ldots, n\} \setminus \{\lambda_1, \ldots, \lambda_d\} \).

In particular:

\[
\forall j \neq j', (w_{j}, w_{j'}) \in \mathfrak{S}_{\lambda_j} \times \mathfrak{S}_{\lambda_{j'}} : w_{j'} w_j = w_j w_{j'} \text{ in } \mathfrak{S}_n.
\]

**Remark 6.2.** If \( w = s_{a_1} \cdots s_{a_r} \in \mathfrak{S}_\lambda \) is a reduced expression, up to a reindexation we know by (6.4) that there is a sequence \( 0 = a_0 \leq a_1 \leq \cdots \leq a_{r-1} \leq a_r = r \) such that for each \( j \in J \), the word \( s_{r-1+1} \cdots s_{r} \) is reduced and lies in \( \mathfrak{S}_{\lambda_j} \). The converse is also true: if for each \( j \in J \) we have a reduced word \( s_{r_j-1+1} \cdots s_{r_j} \in \mathfrak{S}_{\lambda_j} \) then their concatenation \( s_1 \cdots s_r \in \mathfrak{S}_\lambda \) is reduced.
The following proposition is straightforward.

**Proposition 6.3.** The stabiliser of $t^λ$ under the action of $S_n$ is exactly $S_λ$.

We now study the right cosets in $S_n$ for the (left) action of $S_λ$.

**Lemma 6.4.** Two words $w, w' ∈ S_n$ are in the same right coset if and only if $w^{-1} ⋅ t^λ = w'^{-1} ⋅ t^λ$.

The proof is straightforward from Proposition 6.3. An element $C ∈ S_λ \setminus S_n$ is thus determined by the constant value $t := w^{-1} ⋅ t^λ ∈ J^λ$ for $w ∈ C$: we write $C_t$ for the coset $C$ (as each $t ∈ J^n$ has a unique shape, we do not need to precise the underlying composition in the indexation). Noticing that $m_λ = |S_n|/|S_λ|$, we conclude that the cosets are parametrised by the whole set $J^λ$, that is, $S_λ \setminus S_n = \{C_t\}_{t ∈ J^n}$.

We know by [GePf, Proposition 2.1.1] that each coset $C_t$ has a unique minimal length element: we write $π_t ∈ C_t$ for this unique element. In particular, since Lemma 6.4 gives:

$$∀w ∈ S_n, w ∈ C_t ⇔ w ⋅ t = t^λ,$$

we get the following proposition.

**Proposition 6.5.** The element $π_t$ is the unique minimal length element of $S_n$ such that:

$$π_t ⋅ t = t^λ.$$ (6.6)

**Remark 6.6.** The decomposition into right cosets is obtained in the following way. Given $w ∈ S_n$, we know that $w$ belongs to the coset $C_t$ with $t := w^{-1} ⋅ t^λ$. The element $w^{-1} := π_tw^{-1}$ stabilises $t^λ$, thus lies in $S_λ$ and we have $w = wπ_t$.

**Proposition 6.7.** The elements $π_t$ are given by:

$$∀a ∈ \{1, ..., n\}, π_t(a) = λ_{t_{a-1}} + \# \{b ≤ a : t_b = t_a\}.$$ (6.7)

An example is given in Figure 2. To prove Proposition 6.7, we will use the vocabulary of “tableaux”

$$\begin{align*}
\{1, ..., 6\} & \quad 1 \quad 2 \quad 3 \quad 4 \quad 5 \quad 6 \\
t & \quad 3 \quad 1 \quad 3 \quad 2 \quad 3 \quad 1 \\
t^λ & \quad 1 \quad 1 \quad 2 \quad 3 \quad 3 \quad 3 \\
\{1, ..., 6\} & \quad 1 \quad 2 \quad 4 \quad 5 \quad 6 
\end{align*}$$

Figure 2: The permutation $π_t$ for $λ := (2, 1, 3) ≡_3 6$ and $t := (3, 1, 2, 3, 1, 4)$.

(see for example [Ma, §3.1]). As a quick reminder, a $λ$-tableau $T$ is a bijection $\{(j, m) ∈ \mathbb{N}^2 : 1 ≤ j ≤ d \text{ and } 1 ≤ m ≤ λ_j\} → \{1, ..., n\}$; the tableau $T$ is row-standard if in each rows, its entries increase from left to right. Here are two examples of $λ$-tableaux, with $λ := (2, 1, 3) ≡_3 6$:

$$\begin{align*}
T_t & = \begin{array}{c}
2 & 6 \\
4 \\
1 & 3 & 5
\end{array} & 1 & 2 & 3 & 4 & 5 & 6 \\
T_λ & = \begin{array}{c}
1 & 2 \\
3 \\
4 & 6 & 1 & 4
\end{array}
\end{align*}$$

The first only being row-standard.

To any $t ∈ J^λ$, we associate the $λ$-tableau $T_t$ given by the following rule: for $j ∈ J$ and $m ∈ \{1, ..., λ_j\}$, we label the node $(j, m)$ by the index $a$ of the $m$th occurrence of $j$ in $t$, that is, by the integer $a ∈ \{1, ..., n\}$ determined by:

$$t_a = j \text{ and } \# \{b ≤ a : t_b = t_a\} = m.$$ (6.7)

In particular, the tableau $T_t$ is row-standard; conversely, each row-standard $λ$-tableau is a $T_t$ for a unique $t ∈ J^λ$. With the notation of Figure 2, here are two examples of row-standard $λ$-tableaux:

$$\begin{align*}
T_t & = \begin{array}{c}
2 & 6 \\
4 \\
1 & 3 & 5
\end{array} & T_λ & = \begin{array}{c}
1 & 2 \\
3 \\
4 & 6 & 1 & 4
\end{array}
\end{align*}$$

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We consider the natural left action of the symmetric group \( S_n \) on the set of \( \lambda \)-tableaux: if \( w \in S_n \) and \( T \) is a \( \lambda \)-tableau, the tableau \( w \cdot T \) is obtained by applying \( w \) in each box of \( T \). If now \( T \) and \( T' \) are two \( \lambda \)-tableaux, we write \( T \sim T' \) if for all \( j \in J \), the labels of the \( j \)th row of \( T \) are a permutation of the labels of the \( j \)th row of \( T' \).

**Lemma 6.8.** For \( w \in S_n \) and \( t \in J^\lambda \) we have:

\[
w \cdot T_i \sim T_{\chi} \iff w \cdot t = t^\lambda.
\]

**Proof.** If \( j \in J \) and \( m \in \{1, \ldots, \lambda_j\} \), we denote by \( a[j, m] \) the label of the box \((j, m)\) of \( T_i \); by (6.7) we have \( t_{i[j,m]} = j \). We get:

\[
\begin{align*}
w \cdot T_i \sim T_{\chi} & \iff \forall j \in J, \forall m \in \{1, \ldots, \lambda_j\}, w(a[j, m]) \in \{\lambda_j, 1, \ldots, \lambda_j\} \\
& \iff \forall j \in J, \forall m \in \{1, \ldots, \lambda_j\}, t_{w(a[j, m])} = j \\
& \iff \forall j \in J, \forall m \in \{1, \ldots, \lambda_j\}, t_{w(a[j, m])} = t_{a[j, m]} \\
& \iff \forall a \in \{1, \ldots, n\}, t_{\omega(a)} = t_a \\
& \iff \forall a \in \{1, \ldots, n\}, t_a = t_{\omega^{-1}(a)} \\
w \cdot T_i \sim T_{\chi} & \iff t^\lambda = w \cdot t,
\end{align*}
\]

as desired. \( \square \)

**Proof of Proposition 6.7.** Let \( t \in J^\lambda \). There is a unique element \( d(t) \in S_n \) such that \( T_i = d(t) \cdot T_{\chi} \), that is, \( d(t)^{-1} \cdot T_i = T_{\chi} \). By the equation of the coset \( C_i \) given at (6.5) and Lemma 6.8, we get that \( d(t)^{-1} \in C_i \). Applying [Ma, Proposition 3.3], we know that \( d(t)^{-1} \) is the unique minimal length element of \( C_i \). As a consequence, we have \( d(t)^{-1} = \pi_1 \) and thus:

\[
\pi_1 \cdot T_i = T_{\chi}.
\]

Let \( j \in J \) and \( m \in \{1, \ldots, \lambda_j\} \), and let \( a \) (respectively \( \alpha \)) be the label of the box \((j, m)\) in \( T_i \) (resp. \( T_{\chi} \)). In particular, by (6.7) we have \( a = \lambda_j + 1 + m \). Moreover, by (6.8) we have \( \pi_1(a) = \alpha \): we conclude that the announced formula is satisfied, since, by a last use of (6.7), we have \( j = t_a \) and \( m = \# \{ b \leq a \mid t_b = t_a \} \).

\( \square \)

**Remark 6.9.** In [Ma], the author considers the elements of \( S_n \) as acting on \( \{1, \ldots, n\} \) from the right, by \( iw := w(i) \) where \( i \in \{1, \ldots, n\} \) and \( w \in S_n \) is a permutation. This is the right action of \( S_n^\lambda \): in such a setting, we read products of permutations from left to right.

**Lemma 6.10.** Let \( t \in J^\lambda \) and let \( \pi_1 = s_{a_1} \cdots s_{a_r} \) be a reduced expression. Then:

\[
\forall m \in \{1, \ldots, r\}, s_{a_m} \cdot (w_m \cdot t) \neq w_m \cdot t,
\]

where \( w_m := s_{a_{m+1}} \cdots s_{a_r} \) (with \( w_m = 1 \) if \( m = r \)).

**Proof.** Let us suppose \( s_{a_m} \cdot (w_m \cdot t) = w_m \cdot t \) and define \( \tilde{\pi}_1 := s_{a_1} \cdots s_{a_{m-1}} s_{a_{m+1}} \cdots s_{a_r} \). Using the assumption and the equality \( \pi_1 \cdot t = t^\lambda \), we see that the element \( \tilde{\pi}_1 \) verifies \( \tilde{\pi}_1 \cdot t = t^\lambda \) too. As the element \( \tilde{\pi}_1 \) is strictly shorter that \( \pi_1 \) (since \( s_{a_1} \cdots s_{a_r} \) is reduced), this is in contradiction with Proposition 6.5.

\( \square \)

**Remark 6.11.** Using \( t = \pi_1^{-1} \cdot t^\lambda \) in Lemma 6.10, we get the following similar result for \( \pi_1^{-1} \). If \( \pi_1^{-1} = s_{a_r} \cdots s_{a_1} \) is a reduced expression, then:

\[
\forall m \in \{1, \ldots, r\}, w_m' \cdot t^\lambda \neq s_{a_m} \cdot (w_m' \cdot t^\lambda),
\]

where \( w_m' := s_{a_{m+1}} \cdots s_{a_r} \).

The next two lemmas are not essential to the proof of the main theorem of this section, Theorem 6.26; however, they will allow us to relate our construction to the one of [JaPA, PA]. Let \( t \in J^n \) and \( a \in \{1, \ldots, n-1\} \). We give in the next lemma the decomposition of Remark 6.6 for the element \( \pi_1 s_a \); this is in fact a particular case of Deodhar’s lemma (see, for instance, [GeP1, Lemma 2.1.2]).

**Lemma 6.12.** Let \( t \in J^n \) and \( a \in \{1, \ldots, n-1\} \). The element \( \pi_1 s_a \) belongs to the coset \( C_{n-1} \), more precisely we have:

\[
\pi_1 s_a = \begin{cases} 
\pi_1 t_a & \text{if } t_a = t_{a+1}, \\
\pi_{a+1} & \text{if } t_a \neq t_{a+1}.
\end{cases}
\]
Proof. First, from (6.5) we have \( \pi_t s_a \cdot (s_a \cdot t) = t^k \) (where \( \lambda |_{=d} n \) is the shape of \( t \)) thus \( \pi_t s_a \) lies in the coset \( C_{s_a} t \).

We suppose that \( t_a = t_a + 1 \). We have \( \pi_t s_a \pi^{-1}_t = (\pi_t(a), \pi_t(a + 1)) \), and we conclude since \( \pi_1(a + 1) = \pi_1(a) + 1 \) by Proposition 6.7.

We now suppose that \( t_a \neq t_a + 1 \). Using the same Proposition 6.7, we know that the permutation \( w := \pi^{-1}_t \pi_a \in S_n \) is supported in \( \{a, a + 1\} \); thus either \( w = s_a \) or \( w = \text{id} \). Since \( t \neq s_a \cdot t \) we have \( \pi_1 \neq \pi_a \cdot t \), hence \( w \neq \text{id} \). Hence, we get \( w = s_a \), that is, \( \pi_t s_a = \pi_a \cdot t \).

We now generalise the result of Lemma 6.12 in the case \( t_a = t_a + 1 \).

**Lemma 6.13.** Let \( t \in J^n \) and \( a \in \{1, \ldots, n - 1\} \) with \( t_a = t_a + 1 \). Let \( s_b_1 \cdots s_b_r \) be a reduced expression of \( \pi_t \) and set \( w_m := s_{b_{m+1}} \cdots s_{b_r} \). If \( b \in \{1, \ldots, n - 1\} \) verifies \( s_b = w_m s_a w^{-1}_m \) for some \( m \in \{0, \ldots, r\} \), then:

\[
\pi_t s_a = s_b_1 \cdots s_{b_m} s_b s_{b_{m+1}} \cdots s_{b_r},
\]

is a reduced expression. Moreover, every reduced expression of \( \pi_t s_a \) is as above.

**Proof.** We first make an observation. As \( t_a = t_a + 1 \), we deduce from (6.5) that the element \( \pi_t s_a \) remains in \( C_1 \). Hence, by minimality of \( \pi_t \) we have:

\[
\ell(\pi_t s_a) > \ell(\pi_t).
\]

Let now \( s_b_1 \cdots s_{b_r} \) be a reduced expression of \( \pi_t \) and let \( b \in \{1, \ldots, n - 1\} \) and \( m \in \{0, \ldots, r\} \) such that \( s_b = w_m s_a w^{-1}_m \). We have:

\[
\pi_t s_a = s_b_1 \cdots s_{b_m} w_m s_a = s_b_1 \cdots s_{b_m} s_b w_m s_a = s_b_1 \cdots s_{b_m} s_b s_{b_{m+1}} \cdots s_{b_r},
\]

and this expression is reduced since \( \ell(\pi_t s_a) = \ell(\pi_t) + 1 \). Conversely, let \( s_{b'_1} \cdots s_{b'_r} \) be a reduced expression of \( \pi_t s_a \). Since \( \ell((\pi_t s_a) s_a) = \ell(\pi_t s_a) \), we can apply [Hum, §5.8 Theorem]: we know that there is some \( m \in \{0, \ldots, r\} \) such that \( s_{b'_1} \cdots s_{b'_{m+1}} \cdots s_{b'_r} \) is a reduced expression of \( \pi_t \), where the hat denotes the omission. We have:

\[
s_{b'_1} \cdots s_{b'_r} = s_{b'_0} \cdots s_{b'_m} \cdots s_{b'_r} s_a,
\]

thus:

\[
s_{b'_0} \cdots s_{b'_m} = w_m s_a w^{-1}_m,
\]

where \( w_m := s_{b_{m+1}} \cdots s_{b_r} \). We now set \( b := b'_m \) and:

\[
b_p := \begin{cases} 
 b'_{p-1} & \text{if } p \in \{1, \ldots, m\}, \\
 b'_p & \text{if } p \in \{m + 1, \ldots, r\}.
\end{cases}
\]

Moreover:

- the expression \( s_{b_1} \cdots s_{b_r} = s_{b'_0} \cdots s_{b'_{m+1}} \cdots s_{b'_r} = \pi_t \) is reduced;
- we have \( w_m = s_{b'_0} \cdots s_{b'_m} s_{b_{m+1}} \cdots s_{b_r} \);
- we have \( s_b = s_{b'_m} = w_m s_a w^{-1}_m \);

thus the reduced expression \( \pi_t s_a = s_{b'_0} \cdots s_{b'_{m+1}} \cdots s_{b'_r} = s_{b_1} \cdots s_{b_m} s_b s_{b_{m+1}} \cdots s_{b_r} \) is of the desired form.

\[\square\]

**Remark 6.14.** Let \( s_{b_1} \cdots s_{b_r} = \pi_t \) be a reduced expression and set \( w_m := s_{b_{m+1}} \cdots s_{b_r} \). For \( m \in \{0, \ldots, r\} \), there exists \( b' \in \{1, \ldots, n - 1\} \) such that \( s_{b'} = w_m s_a w^{-1}_m \) if and only if \( w_m(a) = w_m(a + 1) \). Moreover, as Lemma 6.10 ensures that \( w_m(a + 1) > w_m(a) \), we have \( w_m(a + 1) = w_m(a) + 1 \) if and only if \( w_m(a + 1) = w_m(a) + 1 \).

We end this subsection by introducing a notation. If \( t \in J^k \) and \( k \in K^3 \), we define:

\[
k^k := \pi^{-1}_t \cdot k \in K^3;
\]

in particular, we may denote by \( k^k \) the elements of \( K^k \).
6.1.3 A “disjoint quiver” Hecke algebra

We consider the setting of §1.2.1. Note that since $\mathcal{K}$ is finite, we can consider the quiver Hecke algebra $H_n(Q)$; recall that its generators are given at (1.29), which are subject to the relations (1.18)–(1.27) and (1.30). We already said that the elements $\psi_n$ do not verify the same braid relations as the elements $s_n \in \mathfrak{S}_n$: in particular, if $s_{n_1} \cdots s_{n_d}$ is another reduced expression for $w$, we may have $\psi_{n_1} \cdots \psi_{n_d} \neq \psi_w$. However, according to Remark 6.2 we can assume that we chose the reduced expressions such that:

$$\forall w = (w_1, \ldots, w_d) \in \mathfrak{S}_\lambda, \psi_w = \psi_{w_1} \cdots \psi_{w_d},$$  

(6.11)

and in the sequel we do suppose that we did so. To that extent, we can first choose some reduced expressions for the elements of the subgroups $\mathfrak{S}_\lambda$ for $j \in J$ and then by product we obtain the reduced expressions of the element of $\mathfrak{S}_\lambda$. Concerning the elements of $\mathfrak{S}_n \setminus \mathfrak{S}_\lambda$, we can choose their reduced expressions arbitrary.

We now suppose that the matrix $Q$ verifies:

$$\forall j \neq j', \forall (k, k') \in K_j \times K_{j'}, Q_{k,k'} = 1.$$  

(6.12)

When the matrix $Q$ is associated with a quiver $\Gamma$ (recall §1.2.2), the condition (6.12) is satisfied when $\Gamma$ is the disjoint union of $d$ proper subquivers $\Gamma^1, \ldots, \Gamma^d$. It means that:

- if $v$ is a vertex in $\Gamma$ then there is a unique $1 \leq j \leq d$ such that $v$ is a vertex of $\Gamma^j$;
- if $(v, w)$ is an edge in $\Gamma$ then there is a (unique) $1 \leq j \leq d$ such that:
  - the vertices $v$ and $w$ are vertices of $\Gamma^j$,
  - the edge $(v, w)$ is an edge of $\Gamma^j$.

Such a disjoint union in $d$ proper subquivers was encountered at (1.38). Moreover, regarding the Cartan matrix of $\Gamma$ we have:

$$\forall j \neq j', \forall (k, k') \in K_j \times K_{j'}, c_{k,k'} = 0;$$  

(6.13)

that is, up to a permutation of the indexing set, the matrix is block diagonal. Finally, for $j \in J$ we define:

$$\forall k, k' \in K_j, Q^{j}_{k,k'} := Q_{k,k'},$$  

(6.14)

in particular for each $j \in J$ and for each $n' \in \mathbb{N}$ we have an associated quiver Hecke algebra $H_n(Q^j)$.

6.1.4 Useful idempotents

We define in this section some idempotents of $H_n(Q)$ which are essential for our proof. Thanks to the defining relations (1.18)–(1.20) and (1.30), for each $\lambda \vdash d \ n$ the following element:

$$e(\lambda) := \sum_{i \in J^\lambda} \sum_{k \in K^i} e(k),$$  

(6.15)

is a central idempotent in $H_n(Q)$, that is, $e(\lambda) = e(\lambda)^2$ commutes with every element of $H_n(Q)$. Moreover:

- if $\lambda' \vdash d \ n$ is different from $\lambda$ then $e(\lambda)e(\lambda') = 0$;
- we have $\sum_{\lambda \vdash d \ n} e(\lambda) = 1$;

hence we have the following decomposition into subalgebras:

$$H_n(Q) = \bigoplus_{\lambda \vdash d \ n} e(\lambda)H_n(Q).$$  

(6.16)

For $t \in J^\lambda$, we also define the following idempotent:

$$e(t) := \sum_{k \in K^i} e(k).$$

We can note that $e(\lambda) = \sum_{i \in J^\lambda} e(t)$. Moreover, we have $e(t)e(t') = 0$ if $t' \in J^n \setminus \{t\}$. We now give some lemmas which involve these elements $e(t)$.
Lemma 6.15. Let $t \in J^\alpha$. We have the following relations:

\[
\psi_a y_{a+1} e(t) = y_a \psi_a e(t) \quad \text{if } t_a \neq t_{a+1},
\]

\[
\psi_a y_{a+1} e(t) = y_a \psi_a e(t) \quad \text{if } t_a \neq t_{a+1},
\]

\[
\psi_a y_{a+1} e(t) = e(t) \quad \text{if } t_a \neq t_{a+1},
\]

\[
\psi_{a+1} \psi_a y_{a+1} e(t) = \psi_a \psi_{a+1} y_{a+1} e(t) \quad \text{if } t_a \neq t_{a+2}.
\]

Proof. Let us first prove the first one. Let $k \in K^1$; we have $k_a \in K_{t_a}$ and $k_{a+1} \in K_{t_{a+1}}$ with $t_a \neq t_{a+1}$ thus $k_a \neq k_{a+1}$. Hence, we get the result using the defining relation (1.24) by summing over all $k \in K^1$. The proofs of the second and the last equalities are similar.

Let us now prove $\psi_d^2 e(t) = e(t)$ if $t_a \neq t_{a+1}$. Let $k \in K^1$; we have $k_a \in K_{t_a}$ and $k_{a+1} \in K_{t_{a+1}}$ with $t_a \neq t_{a+1}$ thus $Q_{k_a, k_{a+1}} = 1$ (see (6.12)). Hence, the defining relation (1.26) gives $\psi_d^2 e(k) = e(k)$, and we again conclude by summing over all $k \in K^1$.

\[\square\]

6.2 About the $\psi_{\pi_1}$

Here we prove some identities which are satisfied by the elements we have just introduced; some of them will be essential to the proof of Theorem 6.26, while others will only be used in §6.4.3, namely with Lemmas 6.20 to 6.22. We first study some properties about the elements $\psi_{\pi_1}$ for $t \in J^\alpha$. We begin by the most important one, which is mentioned in the proof of [SVV, Lemma 3.17].

Lemma 6.16. Let $t \in J^\alpha$. If $s_{a_1} \cdots s_{a_r}$ and $s_{a_1} \cdots s_{b_r}$ are two reduced expressions of $\pi_1$, then:

\[
\psi_{a_1} \cdots \psi_{a_r} e(t) = \psi_{b_1} \cdots \psi_{b_r} e(t).
\]

In other words the element $\psi_{\pi_1} e(t) \in H_n(Q)$ does not depend on the choice of a reduced expression for $\pi_1$.

Proof. By Matsumoto’s theorem, it suffices to check that every braid relation in $s_{a_1} \cdots s_{a_r}$ also occurs in $\psi_{a_1} \cdots \psi_{a_r} e(t)$. By (1.23), it is true for length 2-braids so it remains to check the case of the braids of length 3.

Suppose that we have a braid of length 3 in $s_{a_1} \cdots s_{a_r}$, at rank $m$: we have $a_m = a_{m-2} = a_{m+1} = 1$. We set $a := \min(a_m, a_{m-1})$. With $w_1 := s_{a_1} \cdots s_{a_{a-1}}$ and $w_2 := s_{a_{a+1}} \cdots s_{a_r}$, we have:

\[
w_2(s_a s_{a+1} s_a) w_1 = w_1(s_{a+1} s_a s_{a+1}) w_2,
\]

and we have to prove, with $\psi_1 := \psi_{a_1} \cdots \psi_{a_{a-1}}$ and $\psi_r := \psi_{a_{a+1}} \cdots \psi_{a_r}$:

\[
\psi_1 (\psi_a \psi_{a+1} \psi_a) e(t) = (\psi_{a+1} \psi_a \psi_{a+1}) \psi_r e(t).
\]

Using (1.20), this becomes, where $z := w_r \cdot t$:

\[
\psi_1 (\psi_a \psi_{a+1} \psi_a) e(z) e(t) = (\psi_{a+1} \psi_a \psi_{a+1}) e(z) e(t). \tag{6.17}
\]

By Lemma 6.10, we have $s_{a_{a+1}} \cdot (s_{a_{m+2}} \cdot s) \neq s_{a_{a+2}} \cdot s$. Thus, we have either $s_a \cdot (s_{a+1} \cdot s) \neq s_{a+1} \cdot s$ or $s_{a+1} \cdot (s_a \cdot s) \neq s_{a} \cdot s$; both cases give $s_a \neq s_{a+2}$. Hence, applying Lemma 6.15 we know that (6.17) holds.

Remark 6.17. In particular, if $k \in K^1$ then $\psi_{\pi_1} e(k) \in H_n(Q)$ does not depend on the choice of a reduced expression for $\pi_1$ (note that $\psi_{\pi_1} e(k) = \psi_{\pi_1} e(t) e(k)$).

Similarly to Lemma 6.16, using Remark 6.11 we prove that for $t \in J^\lambda$ the element:

\[
e(t) \psi_{\pi_1} = \psi_{\pi_1} e(t) \in H_n(Q), \tag{6.18}
\]

does not depend on the chosen reduced expression for $\pi_1^{-1}$. We now give some analogues of the results of Lemma 6.15.

Proposition 6.18. Let $t \in J^\lambda$. We have:

\[
\psi_{\pi_1} \psi_{\pi_1} e(t) = e(t),
\]

\[
\psi_{\pi_1} \psi_{\pi_1} e(t^4) = e(t^4).
\]

Remark 6.19. Both factors $\psi_{\pi_1}$ and $\psi_{\pi_1}$ do not depend on the choices of reduced expressions: for instance, using (1.20) we have $\psi_{\pi_1} \psi_{\pi_1} e(t) = \psi_{\pi_1} e(t^4) \psi_{\pi_1}$ thus we can apply Lemma 6.16 and (6.18).
Proof. We only prove the first equality, the proof of the second one being entirely similar. Let \( s_{a_1} \cdots s_{a_r} \) be a reduced expression for \( \pi_t \). We prove by induction that for every \( m \in \{1, \ldots, r+1\} \) we have:

\[
\psi_{\pi_{t}^{m-1}} \psi_{\pi_{t}} e(t) = \psi_{a_r} \cdots \psi_{a_m} \psi_{a_m} \cdots \psi_{a_r} e(t) .
\]  
(6.19)

First, the case \( a = 1 \) comes with the definition of \( \psi_{\pi_{t}^{m-1}} e(t^A) \) and \( \psi_{\pi_{t}} e(t) \). Now, if (6.19) is true for some \( m \in \{1, \ldots, r\} \) we have, using (1.20):

\[
\psi_{\pi_{t}^{m-1}} \psi_{\pi_{t}} e(t) = \psi_{a_r} \cdots \psi_{a_m} e(w_m \cdot t) \psi_{a_m} \cdots \psi_{a_r} ,
\]  
(6.20)

where \( w_m := s_{a_{m+1}} \cdots s_{a_r} \). By Lemma 6.10, we know that \( (w_m \cdot t)_{a_m} \neq (w_m \cdot t)_{a_{m+1}} \). Hence, by Lemma 6.15 we have \( \psi_{\pi_{t}^{m-1}} e(w_m \cdot t) = e(w_m \cdot t) \) thus (6.20) becomes:

\[
\psi_{\pi_{t}^{m-1}} \psi_{\pi_{t}} e(t) = \psi_{a_r} \cdots \psi_{a_m} e(w_m \cdot t) \psi_{a_m} \cdots \psi_{a_r} ,
\]

which becomes, with a last use of (1.20):

\[
\psi_{\pi_{t}^{m-1}} \psi_{\pi_{t}} e(t) = \psi_{a_r} \cdots \psi_{a_{m+1}} \psi_{a_{m+1}} \cdots \psi_{a_r} .
\]

Thus (6.19) holds for every \( m \in \{1, \ldots, r+1\} \), in particular for \( m = r+1 \) we get the statement of the Proposition.

Once again, what follows is not essential to the proof of the main result Theorem 6.26; however, it will allow us to relate our construction to the one of [JaPA, PA]. With a similar proof as Proposition 6.18, we obtain the following lemma.

**Lemma 6.20.** Let \( a \in \{1, \ldots, n\} \) and \( t \in J^\lambda \). We have:

\[
y_{t} \psi_{\pi_{t}} e(t) = \psi_{t} y_{t}^{-1}(a) e(t) ,
\]

\[
y_{t} \psi_{\pi_{t}^{-1}} e(t^A) = \psi_{t}^{-1} y_{t}(a) e(t^A) .
\]

We now want to see what is happening with Lemma 6.12 for the associated elements \( \psi_w \).

**Lemma 6.21.** Let \( t \in J^n \) and \( a \in \{1, \ldots, n-1\} \) such that \( t_a \neq t_{a+1} \). We have:

\[
e(t) \psi_{\pi_{t}^{-1}} \psi_{s_{a+1}} = e(t) \psi_{a} .
\]

**Proof.** By Lemma 6.12 we have \( \ell(\pi_{s_{a+1}}) = \ell(\pi_t) + 1 \). We now simply distinguish cases.

- We first assume that \( \ell(\pi_{s_{a+1}}) = \ell(\pi_t) + 1 \). Hence, applying Lemma 6.16 for the elements \( \psi_{s_{a+1} t} \) and \( \psi_{\pi_t} \), we have \( e(t^A) \psi_{s_{a+1} t} = e(t^A) \psi_{\pi_t} \psi_{a} \). Finally, using Proposition 6.18 we have:

\[
e(t) \psi_{\pi_{t}^{-1}} \psi_{s_{a+1} t} = \psi_{s_{a+1} t} (t) \psi_{\pi_{t}^{-1}} e(t) \psi_{a} = e(t) \psi_{a} .
\]

- We now suppose that \( \ell(\pi_{s_{a+1}}) = \ell(\pi_t) + 1 \); hence, \( \ell(\pi_{t}^{-1}) = \ell(\pi_{s_{a} t}^{-1}) + 1 \). Recall that \( \pi_{t}^{-1} = s_{a} \pi_{s_{a} t}^{-1} \), using the extension (6.18) of Lemma 6.16, we get \( e(t) \psi_{\pi_{t}^{-1}} = e(t) \psi_{s_{a} \pi_{s_{a} t}^{-1}} \) and finally:

\[
e(t) \psi_{\pi_{t}^{-1}} \psi_{s_{a} t} = \psi_{a} e(s_{a} \cdot t) \psi_{s_{a} t}^{-1} \psi_{s_{a} t} = e(t) \psi_{a} .
\]

**Lemma 6.22.** Let \( t \in J^n \) and \( a \in \{1, \ldots, n-1\} \) such that \( t_a = t_{a+1} \). The element \( \psi_{s_{a+1} t} e(t) \in H(a)(Q) \) does not depend on the choice of a reduced expression for \( \pi_{s_{a+1} t} \). In particular:

\[
\psi_{s_{a} t} e(t) = \psi_{s_{a}(a)} \psi_{s_{a+1} t} e(t) .
\]

**Proof.** As in the proof of Lemma 6.16, it suffices to prove that every 3-braid relation which occurs in a reduced expression of \( \pi_{s_{a+1} t} \) is also verified in the corresponding element of \( H(a)(Q)e(t) \). By Lemma 6.13, we know that any reduced expression of \( \pi_{s_{a+1} t} \) can be written \( s_{b_0} \cdots s_{b_r} \), where there is \( m \in \{0, \ldots, r\} \) such that:

- the word \( s_{b_0} \cdots s_{b_m} \) is a reduced expression of \( \pi_t \);

- we have \( s_{b_m} = w_m s_{a} w_m^{-1} \) with \( w_m := s_{b_{m+1}} \cdots s_{b_r} \).
We suppose that a 3-braid appears in $s_{b_0} \cdots s_{b_r}$ at index $l$, that is, we have $b_l = b_{l+2} = b_{l+1} \pm 1$; we set $b := \min(b_l, b_{l+1})$. We want to prove that, with $\psi_1 := \psi_{b_0} \cdots \psi_{b_{l-1}}$ and $\psi_r := \psi_{b_{l+3}} \cdots \psi_{b_r}$:

$$\psi_1(\psi_b \psi_{b+1} \psi_b) \psi_r(e(1)) = \psi_1(\psi_{b+1} \psi_b \psi_{b+1}) \psi_r(e(1)).$$

To that extent, as in the proof of Lemma 6.16, thanks to Lemma 6.15 it suffices to prove that $s_0 \neq s_n+2$ where $s := s_{b+3} \cdots s_{b_r} \cdot t$.

- If $m < l + 1$ then applying Lemma 6.10 we have $s_{b+1} \cdot (s_{b+2} \cdot s) \neq s_{b+2} \cdot s$ thus $s_0 \neq s_0+2$.
- The case $m = l + 1$ is impossible: as $b_l = b_{l+2}$, if $m = l + 1$ we would have $b_{m-1} = b_{m+1}$ and this is nonsense since the expression $s_{b_{m-1}} s_{b_{m+1}}$ is reduced (as a subexpression of the reduced expression $s_{b_1} \cdots s_{b_m} \cdots s_{b_r} = \pi_1$).
- If $m = l + 2$ then by Lemma 6.10 we get $s_0 \cdot (s_{b+1} \cdot s) \neq s_{b} \cdot s$ thus $s_0 \neq s_0+2$.
- Finally, if $m > l + 2$ then we can notice that:

$$s = s_{b+3} \cdots s_{b_r} \cdot t$$

(since $s_{b_m} \cdot (w_m \cdot t) = w_m \cdot (s_{b} \cdot t)$, recall that $s_{b_m} = w_m s_m w_m^{-1}$ and $t_a = t_{a+1}$). Hence, we deduce once again the result from Lemma 6.10.

The last statement of the lemma is now immediate. As $t_a = t_{a+1}$, we can use (6.9) hence $\psi_{\pi_1} e(1) = \psi_{\pi_1} \psi_1 e(1)$. Moreover, applying Lemma 6.12 another consequence of (6.9) is $\ell(\pi_1) = \ell(\pi_2) + 1$, thus we get $\psi_{\pi_1} e(1) = \psi_{\pi_1(\pi_2)} \psi_{\pi_1} e(1)$. Finally, we have $\psi_{\pi_1} e(1) = \psi_{\pi_1(\pi_2)} \psi_{\pi_1(\pi_2)} e(1)$. \hfill $\square$

### 6.3 Decomposition along the subquiver Hecke algebras

We are now ready to prove the main result of this section: we will give in Theorem 6.26 a decomposition of $H_n(Q)$ involving the algebras $H_n(Q')$ for $j \in J$.

#### 6.3.1 A distinguished subalgebra

In this paragraph, we prove the key of Theorem 6.26. Recall that we have set in (6.14):

$$\forall k, k' \in K_j, Q_{k,k'} := Q_{k,k'},$$

for any $j \in J$. Let $\lambda |\sim d n$ be a $d$-composition of $n$. We define the following algebra:

$$H_{\lambda}(Q) := H_{\lambda_1}(Q_1) \otimes \cdots \otimes H_{\lambda_d}(Q_d).$$

With $e_\lambda := e(t^\lambda)$, we prove here that we can identify $H_{\lambda}(Q)$ with the subalgebra $e_\lambda H_n(Q) e_\lambda$ (with unit $e_\lambda$). We reindex the generators $\psi_1, \ldots, \psi_{\lambda_1-1}$ and $y_1, \ldots, y_{\lambda_d}$ of $H_{\lambda_j}(Q_j)$, respectively by $\psi_{\lambda,1}, \ldots, \psi_{\lambda,j-1}$ and $y_{\lambda,j-1+1}, \ldots, y_{\lambda,\lambda_d}$. In particular, we set:

$$\psi^{\otimes} := \psi_{w_1} \otimes \cdots \otimes \psi_{w_d} \in H_\lambda(Q),$$

for $w = (w_1, \ldots, w_d) \in \mathfrak{S}_\lambda$ and:

$$e^{\otimes}(k^1, \ldots, k^d) := e(k^1) \otimes \cdots \otimes e(k^d) \in H_\lambda(Q),$$

for $k = (k^1, \ldots, k^d) \in K_{\lambda_1} \times \cdots \times K_{\lambda_d}$.

**Lemma 6.23.** The following family:

$$\left\{ \psi_{w} y^{r_1}_{1} \cdots y^{r_n}_{n} e(k), w \in \mathfrak{S}_\lambda, r_a \in \mathbb{N}, k \in K^{\lambda_1} \right\},$$

is an $A$-basis of $e_\lambda H_n(Q) e_\lambda$. Moreover, the algebra $e_\lambda H_n(Q) e_\lambda$ is exactly the (non-unitary) subalgebra of $H_n(Q)$ generated by:

- the $\psi_a e_\lambda$ for $a \in \{1, \ldots, n\} \setminus \{\lambda_1, \ldots, \lambda_d\}$;
- the $y_a e_\lambda$ for $1 \leq a \leq n$;
- the $e(k)$ for $k \in K^{\lambda_1}$.

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Proof. The first part is a immediate application of Theorem 1.5, (1.19), (1.20) and Proposition 6.3. It remains to check that \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \) is the described subalgebra, subalgebra that we temporarily denote by \( H \). First, all the listed elements belong to \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \). In particular, for \( a \neq \lambda \), we have indeed \( \psi_a e \lambda \in e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \); we can either use the above basis, or we can simply use (1.20) to get \( \psi_a e \lambda = \psi_a e \lambda \psi_a e \lambda \). Hence, we have \( H \subseteq e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \). Finally, we conclude since every element of the basis (6.22) lies in \( H \).

Lemma 6.24. There is a unitary algebra homomorphism from \( \mathcal{H}_3(Q) \) to \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \).

Proof. We define the algebra homomorphism from \( \mathcal{H}_3(Q) \) to \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \) by sending:

- the generators \( \psi_a^\circ \in \mathcal{H}_3(Q) \) for \( a \in \{1, \ldots, n\} \setminus \{\lambda_1, \ldots, \lambda_d\} \) to \( \psi_a e \lambda \in e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \);
- the generators \( y_a \in \mathcal{H}_3(Q) \) for \( b \in \{1, \ldots, n\} \) to \( y_b e \lambda \in e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \);
- the generators \( e^\circ(k) \in \mathcal{H}_3(Q) \) for \( k \in K \lambda \) to \( e(k) \in e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \).

It suffices now to check the defining relations of \( \mathcal{H}_3(Q) \). We will only check (1.24)–(1.27), the remaining ones being straightforward.

(1.24). Let \( a \notin \{\lambda_1, \ldots, \lambda_d\} \) and \( k \in K \lambda \). If \( k_a \in K_j \), as \( a \neq \lambda_j \) for any \( j \)' we have \( k_{a+1} \in K_j \) (cf. (6.3)). Hence, in \( H_n(Q) \) the relation (1.24):

\[
\psi_a^\circ y_{a+1} e^\circ(k) = \begin{cases} 
(y_a \psi_a^\circ + 1) e^\circ(k) & \text{if } k_a = k_{a+1}, \\
(y_a \psi_a^\circ) e^\circ(k) & \text{if } k_a \neq k_{a+1},
\end{cases}
\]

comes from the corresponding relation in \( \mathcal{H}_n(Q) \). The same relation:

\[
\psi_a y_{a} e^\circ(k) = \begin{cases} 
(y_a \psi_a + 1) e(k) & \text{if } k_a = k_{a+1}, \\
(y_a \psi_a) e(k) & \text{if } k_a \neq k_{a+1},
\end{cases}
\]

is verified in \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \), as a relation in \( H_n(Q) \).

(1.25). Similar.

(1.26). Similarly, the indices \( k_a, k_{a+1} \) are in a same \( K_j \). Hence, the relation (1.26) in \( \mathcal{H}_3(Q) \) comes from a relation in \( \mathcal{H}_n(Q) \), and the same relation is verified in \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \).

(1.27). For \( j \in J \) and \( a \in \{\lambda_1, \ldots, \lambda_j - 1, \ldots, \lambda_j - 2\} \), the relation (1.27) is a relation from \( \mathcal{H}_n(Q) \), and this same relation is verified in \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \).

Proposition 6.25. The previous algebra homomorphism \( \mathcal{H}_3(Q) \to e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \) is an isomorphism. In particular, we can identify \( \mathcal{H}_3(Q) \) to a (non-unitary) subalgebra of \( H_n(Q) \).

Proof. We know by Theorem 1.5 that \( \mathcal{H}_3(Q) \) has for basis:

\[
\{ \psi_{w_1} y_{\lambda_{j-1}+1} \cdots y_{\lambda_j} e(k') : w_j \in \mathcal{S}_\lambda, r_j \in \mathbb{N}, k' \in K_j \lambda \}.
\]

hence, the following family:

\[
\{ \psi_{w_1} y_{\lambda_{j-1}+1} \cdots y_{\lambda_j} e^\circ(k) : w \in \mathcal{S}_\lambda, r \in \mathbb{N}, k \in K_{\lambda_{j-1}} \times \cdots \times K_{\lambda_j} \},
\]

is a basis of \( \mathcal{H}_3(Q) \). We conclude since by (6.11) the homomorphism of Lemma 6.24 sends this basis onto the basis of \( e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \) given in Lemma 6.23. (In particular, note that \( \psi_w \in \mathcal{H}_3(Q) \) for \( w \in \mathcal{S}_\lambda \) is sent to \( \psi_w e \lambda \in e \lambda \Lambda \mathcal{H}_3(Q) e \lambda \).)

6.3.2 Decomposition theorem

We recall the notation \( m_{\lambda} \) for \( \lambda \mid d \) \( n \) introduced at (6.2).

Theorem 6.26. We have an \( A \)-algebra isomorphism:

\[
H_n(Q) \simeq \bigoplus_{\lambda \mid d \mid n} \text{Mat}_{m_{\lambda}} \mathcal{H}_3(Q).
\]
The remaining part of this paragraph is devoted to the proof of Theorem 6.26. Due to (6.16), it suffices to prove that we have an $A$-algebra isomorphism:
\[ e(\lambda)H_n(Q) \simeq \text{Mat}_{m_{\lambda}}H_{\lambda}(Q). \]  
(6.23)

Let us label the rows and the columns of the elements of $\text{Mat}_{m_{\lambda}}H_{\lambda}(Q)$ by $(t',t) \in (J^\lambda)^2$, and let us write $E_{t',t}$ for the elementary matrix with one 1 at position $(t',t)$ and 0 everywhere else. Recall the following property verified by the $E_{t',t}$:
\[ \forall t, t', s, s' \in J^\lambda, E_{t',t}E_{t',s} = \delta_{t,s'}E_{t',s}. \]  
(6.24)

We have the following $A$-module isomorphism, where $t, t' \in J^\lambda$:
\[ e(t')H_n(Q)e(t) \simeq H_\lambda(Q)E_{t',t}. \]
Indeed, let us define:
\[ \Phi_{t',t} : H_\lambda(Q)E_{t',t} \to e(t')H_n(Q)e(t), \]
\[ \Psi_{t',t} : e(t')H_n(Q)e(t) \to H_\lambda(Q)E_{t',t}, \]
(6.25)

by:
\[ \forall v \in H_\lambda(Q), \Phi_{t',t}(vE_{t',t}) := \psi_{v_{t',t}}, \]
\[ \forall w \in e(t')H_n(Q)e(t), \Psi_{t',t}(w) := (\psi_{w_{t',t}}, w^\pi_{t',t}E_{t',t}). \]

The goal sets of (6.25) are respected, according to (1.20), (6.6) and Proposition 6.25. Indeed, for instance we have, for $v \in H_\lambda(Q) \simeq e_3H_n(Q)e_3$:
\[ \Phi_{t',t}(vE_{t',t}) = \psi_{v_{t',t}}^\pi_{t',t} \psi_{v_{t',t}} = \psi_{\pi_{t',t}}^\pi_{t',t} e(t^\lambda) \psi_{t',t} = e(\pi_{t',t}^{-1} \cdot t^\lambda) \psi_{\pi_{t',t}} \psi_{t',t} \psi_{e(t^\lambda)}, \]
\[ \Phi_{t',t}(vE_{t',t}) = e(t^\lambda) \psi_{\pi_{t',t}}^{-1} \psi_{t',t} \psi_{e(t^\lambda)}(t) \in e(t^\lambda)H_n(Q)e(t). \]

Remark 6.27. Our map $\Phi_{t',t}$ is similar to $[SVV, (20)]$.

Furthermore, these two maps $\Phi_{t',t}$ and $\Psi_{t',t}$ are clearly $A$-linear and by Proposition 6.18 these are inverse isomorphisms. We now set:
\[ \Phi_\lambda := \bigoplus_{t, t' \in J^\lambda} \Phi_{t',t} : \text{Mat}_{m_\lambda}H_\lambda(Q) \to e(\lambda)H_n(Q), \]
\[ \Psi_\lambda := \bigoplus_{t, t' \in J^\lambda} \Psi_{t',t} : e(\lambda)H_n(Q) \to \text{Mat}_{m_\lambda}H_\lambda(Q). \]
(6.26)

From the properties of $\Phi_{t',t}$ and $\Psi_{t',t}$, the above maps are inverse $A$-module isomorphisms; it now suffices to check that $\Phi_\lambda$ is an $A$-algebra homomorphism. This property comes from the following one:
\[ \Psi_{t',t}(w_{t',t})\Psi_{s',s}(w_{s',s}) = \Psi_{t',t}(w_{t',t}w_{s',s}), \]
(6.27)

where $t, t', s, s' \in J^\lambda$, $w_{t',t} \in e(t')H_n(Q)e(t)$ and $w_{s',s} \in e(s')H_n(Q)e(s)$. The equality (6.27) is obviously satisfied when $t \neq s'$ since both sides are zero, thus we assume $t = s'$. We have, using Proposition 6.18 and noticing that $w_{t,s} = e(t)w_{t,s}$:
\[ \Psi_{t',t}(w_{t',t})\Psi_{s',s}(w_{s',s}) = \Psi_{t',t}(w_{t',t}w_{s',s}) = (\psi_{t',t}^\pi_{t',t})^\pi_{t',t}(\psi_{t',t}w_{t,s}^\pi_{t',t})E_{t',t}E_{t,s} = \psi_{t',t}^\pi_{t',t}w_{t,s}(t^\lambda)w_{t,s}^\pi_{t',t}E_{t',t}E_{t,s} = \psi_{t',t}^\pi_{t',t}w_{t,s}^\pi_{t',t}E_{t',t}E_{t,s} = \Psi_{t',t}(w_{t',t}w_{s',s}). \]

Finally, the maps $\Phi_\lambda$ and $\Psi_\lambda$ are inverse $A$-algebra isomorphisms; we deduce the isomorphism (6.23) and thus Theorem 6.26.
Remark 6.28. For $k \in \mathcal{K}^\lambda$, we write here $k^* := \pi_1, k \in \mathcal{K}^\lambda$. Using Proposition 6.18 and Lemmas 6.20, 6.22, we can give the images of the generators of $e(\lambda)H_n(Q)$ for each $t \in J^\lambda$ and $k \in \mathcal{K}^\lambda$:

$$\Psi_\lambda(e(k)) = e(k^*)E_{t; t},$$

for $a \in \{1, \ldots, n\}$, $\Psi_\lambda(y_{a}e(k)) = y_{e(a)}e(k^*)E_{t; t}$,

$$\forall a \in \{1, \ldots, n - 1\}, \Psi_\lambda(y_{a}e(k)) = \begin{cases} e(k^*)E_{a; t} & \text{if } t_a \neq t_a + 1, \\ \psi_{e(a)}(k^*)E_{t; t} & \text{if } t_a = t_a + 1 \end{cases}$$

(note that $\pi_{s_a+1}^{-1} = \begin{cases} \text{id} & \text{if } t_a \neq t_a + 1, \\ \pi_{e(a)} & \text{if } t_a = t_a + 1 \end{cases}$, cf. Lemma 6.12). We observe that these images look like those described in [JaPA, (22)] and [PA, (3.2)–(3.4)].

Remark 6.29. We consider the setting of Proposition 1.7. We can prove that the algebra isomorphism $H_n(Q) \cong \oplus_{\lambda \vdash d, \emptyset} \text{Mat}_{m_\lambda} \Lambda(Q)$ is a graded isomorphism (with the canonical gradings on the direct sum, matrix algebras and tensor products). In particular, we have (recall the notation $k^1$ of (6.10)):

$$\deg \psi_{e(a)}(k) = \deg \psi_{e(a)}(k^1) = 0,$$

as a consequence of Lemma 6.10 and Remark 6.11. Indeed, if for $s \in J^n$ and $a \in \{1, \ldots, n - 1\}$ we have $s_a \neq s_{a + 1}$ then for any $k \in \mathcal{K}^\lambda$ we have $c_{s_a, s_{a + 1}} = 0$ (cf. (6.13)).

6.4 Cyclotomic version

Let us consider a weight $\Lambda = (\Lambda_k)_{k \in \mathbb{N}} \in \mathcal{N}^\mathbb{C}$; for $j \in J$, we write $\Lambda^j \in \mathcal{N}^{\mathbb{C}^j}$ the restriction of $\Lambda$ to $\mathcal{K}_j$. We show how the isomorphism of Theorem 6.26 is compatible with cyclotomic quotients, as defined in (1.33).

6.4.1 Factorisation theorem

For $\lambda \vdash d, n$, we define the cyclotomic quotient of $H_\lambda(Q)$ by:

$$H^\lambda_n(Q) := H^\lambda_n(Q^1) \otimes \cdots \otimes H^\lambda_n(Q^d),$$

that is, $H^\lambda_n(Q)$ is the quotient of $H_\lambda (Q)$ by the two-sided ideal generated by the elements:

$$\sum_{m=0}^{\lambda_k} c_m y_{a}^m e(k) = 0 \quad \forall k \in \mathcal{K}^\lambda, \forall a \in \{\lambda_0 + 1, \ldots, \lambda_{d - 1} + 1\}. \quad (6.28)$$

Theorem 6.30. The isomorphism of Theorem 6.26 factors through the cyclotomic quotients, in other words we have:

$$H^\lambda_n(Q) \cong \bigoplus_{\lambda \vdash d, \emptyset} \text{Mat}_{m_\lambda} H^\lambda_n(\Lambda).$$

Proof. Let $\lambda \vdash d, n$ and let $\mathcal{J}$ (resp. $\mathcal{J}_\lambda$) be the two-sided ideal of $e(\lambda)H_n(Q)$ (resp. $H_\lambda(Q)$) generated by the elements in (1.33) (resp. (6.28)). It suffices to prove that $\Psi_\lambda(\mathcal{J}) = \text{Mat}_{m_\lambda} \mathcal{J}_\lambda$: we will prove $\Psi_\lambda(\mathcal{J}) \subseteq \text{Mat}_{m_\lambda} \mathcal{J}_\lambda$ and $\mathcal{J} \supseteq \Phi_\lambda(\text{Mat}_{m_\lambda} \mathcal{J}_\lambda)$.

- Each element of $\mathcal{J}$ can be written as:

$$\sum_{\nu, \omega \in H_n(Q)} \sum_{\lambda \vdash d} \sum_{k \in \mathcal{K}^\lambda} \sum_{m=0}^{\lambda_k} c_m y_{\nu}^m e(k^1) \nu.$$  

By (1.20), Proposition 6.18 and Lemma 6.20, the previous element becomes:

$$\sum_{\nu, \omega \in H_n(Q)} \sum_{\lambda \vdash d} \sum_{k \in \mathcal{K}^\lambda} \sum_{m=0}^{\lambda_k} \psi_{\nu} c_m y_{\omega}^m e(\pi_k \cdot k^1) \psi_{\nu}^* \omega.$$
As $\Psi_\lambda$ is an algebra homomorphism and $\text{Mat}_{m,\lambda} J_\lambda$ is a two-sided ideal, it suffices to prove that for any $k \in K^\lambda$ and $t \in J^\lambda$ we have, with $a := \pi_\lambda(1)$:

$$\sum_{m=0}^{\Lambda_{k,\lambda}} c_m \Psi_\lambda(y_m^k e(k)) \in \text{Mat}_{m,\lambda}(J_\lambda).$$

The above element is exactly:

$$\sum_{m=0}^{\Lambda_{k,\lambda}} c_m y_m^k e(k) E_{t',t}$$

(recall that $\pi_\lambda = \text{id}$), hence we are done since the element $\sum_{m=0}^{\Lambda_{k,\lambda}} c_m y_m^k e(k)$ lies in $J_\lambda$; in particular, there is a $j \in J$ such that $a = \lambda_{j-1} + 1$, cf. Proposition 6.7.

- Each element of $\text{Mat}_{m,\lambda} J_\lambda$ can be written:

$$\sum_{t,t' \in J^\lambda} \sum_{v,w \in H_\lambda(Q)} \sum_{k \in K^\lambda} t_{\lambda,\lambda}^{a \in \{\lambda_0 + 1, \ldots, \lambda_{d-1} + 1\}} \sum_{a \in \{\lambda_0 + 1, \ldots, \lambda_{d-1} + 1\}} \sum_{m=0}^{\Lambda_{k,\lambda}} c_m y_m^k e(k) E_{t',s}.$$

As $\Phi_\lambda$ is an algebra homomorphism and $J$ is a two-sided ideal, it suffices to prove that for any $t, t' \in J^\lambda, k \in K^\lambda$ and $a = \lambda_{j-1} + 1$ for $j \in J$ we have:

$$\sum_{m=0}^{\Lambda_{k,\lambda}} c_m \Phi_\lambda(y_m^k e(k) E_{t',s}) \in J.$$

We consider an element $s \in J^\lambda$ which verifies $s_1 = j$ (we can take for instance $s := (1, a) \cdot t^\lambda$); note that $\pi_\lambda(1) = \lambda_{j-1} + 1 = a$. Using (6.24), we can write the above element:

$$\sum_{m=0}^{\Lambda_{k,\lambda}} c_m \Phi_\lambda(y_m^k e(k) E_{t',s}) \Phi_\lambda(E_{s,\lambda}),$$

hence it suffices to prove that

$$\alpha := \sum_{m=0}^{\Lambda_{k,\lambda}} c_m \Phi_\lambda(y_m^k e(k) E_{t',s}),$$

belongs to the ideal $J$. We get:

$$\alpha = \sum_{m=0}^{\Lambda_{k,\lambda}} c_m \psi_{n-1}^{a-1} y_m^k e(k) \psi_\lambda s$$

$$= \sum_{m=0}^{\Lambda_{k,\lambda}} c_m \psi_{n-1}^{a-1} \psi_\lambda s y_m^k e(k)$$

$$= \psi_{n-1}^{a-1} \psi_\lambda s \sum_{m=0}^{\Lambda_{k,\lambda}} c_m y_m^k e(k^a),$$

where we recall the notation $k^a \in K^a$ from (6.10). We have $k^a_n = k_n$, thus we get:

$$\alpha = \psi_{n-1}^{a-1} \psi_\lambda s \sum_{m=0}^{\Lambda_{k,\lambda}} c_m y_m^k e(k^a),$$

and we are done.
6.4.2 An alternative proof

We explain here how we can get Theorem 6.30 from [SVV, §3.2.3]. We make the following assumption:

\[ Q_{k,k'}(u,v) = r_{k,k'}(u-v)^{-c_{k,k'}}, \]

where the \( r_{k,k'}, c_{k,k'} \) are some scalars; we refer to [SVV, §3.2.2] for more details about the setting. In particular, we suppose that \( \Lambda = \sum_{k \in K} \Lambda_k \omega_k \) where the \( \omega_k \) are the fundamental weights related to the ambient Cartan datum. We define the set \( K_\Lambda \) by (cf. [SVV, §3.1.3]):

\[ K_\Lambda := \{(k, m) : k \in K, m \in \{1, \ldots, \Lambda_k\}\}. \]

we will also use the sets \( K_{\Lambda, j} := (K_j)_{\Lambda, j} \) for \( j \in J \). For \( t = (k, m) \in K_\Lambda \), we write \( k_t := k \) and \( \omega_t := \omega_k \). Finally, for convenience we introduce some notation:

- we write \( |\cdot|_\Lambda \) instead of \( |\cdot|_{K_\Lambda} \);
- if \( \mu \) (respectively \( \nu \)) is an \( r \)-composition \( s \)-composition, we set \( m_{\mu} := \frac{\nu_1 \cdots \nu_r}{\mu_1 \cdots \mu_s} \); in particular with the 1-partition \( \nu := (\mu_1 + \cdots + \mu_r) \) we recover \( m_{\nu} = m_\mu \).

**Theorem 6.31** ([SVV, Theorem 3.15]). There is an algebra isomorphism

\[ H^n_\Lambda(Q) \cong \bigoplus_{(n_t) \in K_\Lambda} \text{Mat}_{m_{\nu}} \left( \bigotimes_{t \in K_\Lambda} H^n_{\nu_t}(Q) \right). \]

If we look closer to the proof of [SVV], comparing to the proof of Theorem 6.26 the idea is still to consider some elements \( e(t) \) but with idempotents which "refine" the idempotents \( e(k) \) (these idempotents are indexed by \( \Lambda_\Lambda \), which is much bigger than \( K^\Lambda \)). In particular, the following isomorphism for any \( \lambda = d \) \( n \), writing \( (n_t)_t \) for the restriction of \( (n_t)_t \in K_\Lambda \) to \( K_{\Lambda, j} \):

\[ \text{Mat}_{m_{\nu}} H^n_\Lambda(Q) \cong \bigoplus_{(n_t)_t \in K_{\Lambda, j}} \text{Mat}_{m_{\nu}} \left( \bigotimes_{t \in K_{\Lambda, j}} H^n_{\nu_t}(Q) \right) \] (6.29)

implies our Theorem 6.30 by summing over all \( \lambda = d \) \( n \). In order to prove (6.29), we can simply apply Theorem 6.31 to the factors \( H^n_\Lambda(Q)^j \) of \( H^n_\Lambda(Q) \). We have, for \( j \in J \):

\[ H^n_\Lambda(Q^j) \cong \bigoplus_{(n_t)_t \in K_{\Lambda, j}} \text{Mat}_{m_{\nu}} \left( \bigotimes_{t \in K_{\Lambda, j}} H^n_{\nu_t}(Q^j) \right). \] (6.30)

Before going further, we give the following lemma. Let us mention that we can find a non-cyclotomic statement in [Rou, Corollary 3.8]; see also [Ma2, Proposition 2.4.6] and [BoMa, Lemma 1.16].

**Lemma 6.32.** Let \( j \in J \). If \( k \in K_j \) then \( H^n_k(Q) \cong H^n_k(Q^j) \).

**Proof.** It suffices to prove that every \( e(k) \in H_n(Q) \) with \( k \in K^\Lambda \) \( K^\Lambda \) vanishes in \( H^n_k(Q) \). To that extent, we prove by induction on \( a \in \{1, \ldots, n\} \) the following statement:

\[ \forall k \in K^\Lambda, \exists b \in \{1, \ldots, a\}, k_b \not\in K_j \implies e(k) = 0 \text{ in } H^n_k(Q). \] (6.31)

First, we shall verify this proposition for \( a = 1 \): let \( k \in K^\Lambda \) such that \( \exists b \in \{1, \ldots, 1\}, k_b \not\in K_j \). We obviously have \( k_1 \not\in K_j \), in particular \( k_1 \neq k \) thus it follows directly from the cyclotomic condition (1.33) that \( e(k) = 0 \) in \( H^n_k(Q) \). We now assume that (6.31) is verified for some \( a \in \{1, \ldots, n - 1\} \) and we let \( k \in K^\Lambda \) such that \( \exists b \in \{1, \ldots, a + 1\}, k_b \not\in K_j \). We know from the induction hypothesis that \( e(k) = 0 \) in \( H^n_k(Q) \) if \( k_a \not\in K_j \) or \( k_{a + 1} \in K_j \), hence it remains to deal with the case \( k_a \in K_j \) and \( k_{a + 1} \not\in K_j \). Recalling (6.12), this implies that \( Q_{k_a, k_{a + 1}} = 1 \). In particular, the defining relation (1.26) becomes:

\[ \psi^2 e(k) = e(k). \] (6.32)

Besides, using (1.20) and the induction hypothesis, we have:

\[ \psi_a e(k) = e(k) \psi_a = 0 \text{ in } H^n_k(Q). \]

Thus, left-multiplying by \( \psi_a \) and using (6.32) we get \( e(k) = 0 \) in \( H^n_k(Q) \) which ends the induction. Finally, for \( a = n \), we get that if \( k \in K^\Lambda \) then \( e(k) = 0 \) in \( H^n_k(Q) \). \( \square \)
We now go back to the proof of (6.29). Using Lemma 6.32, the isomorphism (6.30) gives:

\[
\text{H}^k(Q) \simeq \bigotimes_{\lambda \in \mathcal{K}} \text{H}^{\nu}_\lambda(Q) = \bigotimes_{\lambda \in \mathcal{K}} \bigoplus \text{Mat}_m(\nu_\lambda) \left( \bigotimes_{\ell \in \mathcal{K}_A} \text{H}^{\nu}_\ell(Q) \right).
\]

We obtain:

\[
\text{H}^k(Q) \simeq \bigoplus_{\lambda \in \mathcal{K}} \bigoplus_{\ell \in \mathcal{K}_A} \text{Mat}_m(\nu_\lambda) \left( \bigotimes_{\ell \in \mathcal{K}_A} \text{H}^{\nu}_\ell(Q) \right).
\]

Finally, we deduce (6.29) and thus Theorem 6.30 from the equality \(m_{\lambda m_{(n_1)}} = m_{(n_1)}\).

### 6.4.3 An application

Here we explain how our Theorem 6.30 is related to our previous result Theorem 4.1. We suppose that \(A = \mathbb{C}\) and that all the \(\mathcal{K}_j\) for \(j \in I = \{1, \ldots, d\}\) have the same finite cardinality \(e \geq 2\) for each of these parts. Finally, we consider \(q \in \mathbb{C}^\times\) a primitive \(e\)th root of unity, and we write:

\[
\text{BK} : \tilde{\text{H}}^k(Q) \simeq \text{H}^k(\Gamma_e)
\]

for the \(\mathbb{C}\)-algebra isomorphism of [BrKl].

We have an algebra isomorphism, constructed by Poulain d’Andecy [PA]:

\[
\text{JPA} : \tilde{\text{S}}^\Lambda_m(Q) \simeq \bigotimes_{\lambda \in \mathcal{K}} \text{Mat}_m(\tilde{\text{H}}^k(Q)),
\]

where \(\tilde{\text{H}}^k(Q) := \tilde{\text{H}}^k_\lambda(Q) \otimes \cdots \otimes \tilde{\text{H}}^k_\lambda(Q)\). This isomorphism is a generalisation of the main result of [JaPA] (which is in fact a particular case of a result of Lusztig [Lu]), and is defined on the generators as follows:

\[
\text{JPA}(t_a) = \sum_{\ell \in J^m} \xi_{\lambda \ell} E_{1,\ell}, \quad \text{JPA}(X_a) = \sum_{\ell \in J^m} X_{\ell a} E_{1,\ell}, \quad \text{JPA}(g_a) = \sum_{\ell \in J^m} g_{\ell a} E_{1,\ell} + \sum_{\ell \in J^m} \sqrt{q} E_{1,\ell - 1}.
\]

**Remark 6.33.** Note two slight differences with [PA]:

- our elements \(E_{1,\ell}\) are written \(E_{\chi,\ell}\), where \(\chi\) is a character of \((\mathbb{Z}/d\mathbb{Z})^m = J^m\);
- Poulain d’Andecy considers left cosets instead of our right ones, in particular his minimal length representatives \(\pi_\ell\) verify \(\pi_\ell = e\).

We recall from (1.38) that \(\hat{\Gamma} = \bigcup_{j \in J} \Gamma_e\); in particular, its vertex set is exactly \(\hat{K} \simeq K = I \times J\). The two previous results, together with our Theorem 6.30, gives straightforwardly the following theorem.

**Theorem 6.34.** We have an algebra isomorphism:

\[
\Phi^\Lambda_m \circ \text{BK} \circ \text{JPA} : \tilde{\text{S}}^\Lambda_m(Q) \simeq \text{H}^k(Q),
\]

where:

\[
\Phi^\Lambda_m : \text{H}^k(\Gamma) \rightarrow \text{H}^k(Q).
\]
• the homomorphism \( \Phi_n^\lambda : \oplus_\lambda \text{Mat}_{m,n} \overline{H}_n^\lambda(\Gamma) \rightarrow \text{H}_n^\lambda(\Gamma) \) is the isomorphism of Theorem 6.30, that is, induced by the homomorphism \( \Phi_n : \oplus_\lambda \text{Mat}_{m,n} \text{H}_\lambda(\Gamma) \rightarrow \text{H}_n(\Gamma) \).

An algebra isomorphism \( \overline{\lambda} \) of \( (\Gamma) \) was already constructed in Theorem 6.30: we shall denote it by \( \overline{\lambda} \). An interesting question is to know whether we recover the same isomorphism as above. In other words, does the diagram of Figure 3 commute?

![Figure 3: A commutative diagram?](image)

As we deal with algebra homomorphisms, it suffices to check that the images of the generators of \( \overline{\lambda} \) are the same. We will use the following notation: for \( t \in J^n \) we set \( t' := \pi_t \cdot t \). (With \( \lambda \equiv [t] \), we have of course \( t'' = t^\lambda \).) Moreover, we will keep on using the notation \( t \in J^n \) of Section 6 for the elements we denoted by \( t \in J^n \) from Section 1 to Section 5.

**Image of \( I_n \).** Let \( 1 \leq a \leq n \). Recall from §3.1 that:

\[
\overline{\lambda}(I_n) = \sum_{j \in J^n} e(j) \xi^a = \sum_{i \in J^n} e(i) \xi^a_i \in \text{H}_n^\lambda(\Gamma).
\]

(6.36)

Recalling (6.33), we obtain:

\[
\overline{\lambda}(I_n) = \sum_{i \in J^n} \sum_{k \in K^i} \xi^a_i \epsilon(k) E_{i,k} \in \bigoplus_{\lambda | d^n} \text{Mat}_{m,n} \text{H}_n^\lambda(\Gamma).
\]

Hence, with the usual manipulations, we get:

\[
\Phi_n^\lambda \circ \overline{\lambda}(I_n) = \sum_{i \in J^n} \sum_{k \in K^i} \psi_a^{-1} \psi_1 e_i(k) \epsilon(k)
\]

\[
= \sum_{i \in J^n} \sum_{k \in K^i} \xi^a_i \epsilon(k)
\]

thus it coincides with (6.36).

**Image of \( X_{a^o} \).** (It is in fact enough to study the case \( a = 1 \).) Let \( 1 \leq a \leq n \). Recall from §3.1 that:

\[
\overline{\lambda}(X_{a^o}) = \sum_{i \in J^n} q^{a^o}(1 - y_a) e(i) = \sum_{k \in K^i} q^{a^o}(1 - y_a) \epsilon(k) \in \text{H}_n^\lambda(\Gamma),
\]

(6.37)

where we write \( q^k := q^k \). Let \( (i, j) \in K = I \times J \). Recalling (6.34), we get:

\[
\overline{\lambda}(X_{a^o}(i)) = \sum_{i \in J^n} X_{a^o}(i) E_{i,1} \in \bigoplus_{\lambda | d^n} \text{Mat}_{m,n} \text{H}_n^\lambda(\Gamma).
\]

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We write, where $j := t_a$ and $\lambda := [t]$:

$$X_{\pi_1(a)} = \sum_{k' \in K_{j, a}^{\lambda}} q^{k_{\pi_1(a)}}(1 - y_{\pi_1(a)})e(k') \in H_{j, a}^{\lambda}(\Gamma) \subseteq H_{j}^{\lambda}(\Gamma),$$

where $k' \in K_{j}^{\lambda}$ is indexed by $(\lambda_{j-1} + 1, \ldots, \lambda_{j})$. Hence, we have:

$$X_{\pi_1(a)} = \sum_{k \in K_{\lambda}^{a}} q^{k_{\pi_1(a)}}(1 - y_{\pi_1(a)})e(k) \in H_{\lambda, a}^{\lambda}(\Gamma),$$

and thus:

$$\Phi_{\lambda}^{a} \circ BK \circ JPA(X_{\lambda}) = \sum_{i \in J_\lambda} \sum_{k' \in K_{j, a}^{\lambda}} q^{k_{\pi_1(a)}}(1 - y_{\pi_1(a)})e(k)E_{i,1} = \sum_{i \in J_\lambda} \sum_{k \in K_{\lambda}^{a}} q^{k_{\pi_1(a)}}(1 - y_{\pi_1(a)})e(k)\psi_{\pi_1} = \sum_{i \in J_\lambda} \sum_{k' \in K_{j, a}^{\lambda}} q^{k_{\pi_1(a)}}(1 - y_{\pi_1(a)})e(k') \in H_{\lambda, a}^{\lambda}(\Gamma),$$

which is (6.37).

**Image of $g_a$.** Let $1 \leq a < n$. Recall from §3.1 that:

$$\text{BK}(g_a) = \sum_{k \in K_{\lambda}^{a}} (\psi_{\pi_1(a)}Q_a(k) - P_a(k))e(k) \in H_{\lambda}^{a}(\Gamma),$$

where $Q_a(k), P_a(k) \in \mathbb{C}[[y_a, y_{a+1}]]$ are some power series. For convenience, we write $\widehat{Q}_a(k), \widehat{P}_a(k) \in \mathbb{C}[[Y_a, Y_{a+1}]]$ the underlying power series, which verify $\widehat{Q}_a(k)(y_a, y_{a+1}) = Q_a(k)$ and $\widehat{P}_a(k)(y_a, y_{a+1}) = P_a(k)$. These power series depend only on $k_a$ and $k_{a+1}$; that is, $\widehat{Q}_a(k)(Y, Y') = \widehat{Q}(k')(Y, Y')$ and $\widehat{P}_a(k)(Y, Y') = \widehat{P}(k')(Y, Y')$ if $k_a = k'_a$ and $k_{a+1} = k'_{a+1}$. Moreover, recall that if $k$ and $s_a \cdot k$ are labellings of two different $t \in J^n$, we have $\widehat{Q}_a(k) = \sqrt{q}$ and $\widehat{P}_a(k) = 0$ (cf. Remark 2.11).

Recalling (6.35), we get:

$$\text{BK}(g_a) = \sum_{i \in J_\lambda} \sum_{k \in K_{\lambda}^{a}} g_{\pi_1(a)}(E_{i,1}) + \sum_{i \in J_\lambda} \sum_{k \in K_{\lambda}^{a}} \sqrt{q}E_{i,1, a} \in \bigoplus_{\lambda \in \mathbb{n}} \text{Mat}_{\lambda, \lambda} H_{\lambda}^{\lambda}(\Gamma).$$

With $j := t_a$ and $\lambda := [t]$, we have:

$$g_{\pi_1(a)} = \sum_{k' \in K_{j, a}^{\lambda}} (\psi_{\pi_1(a)}Q_{\pi_1(a)}(k') - P_{\pi_1(a)}(k'))e(k') \in H_{j, a}^{\lambda}(\Gamma) \subseteq H_{j}^{\lambda}(\Gamma)$$

(recall that $k' \in K_{j}^{\lambda}$ is indexed by $(\lambda_{j-1} + 1, \ldots, \lambda_{j})$), hence:

$$g_{\pi_1(a)} = \sum_{k \in K_{\lambda}^{a}} (\psi_{\pi_1(a)}Q_{\pi_1(a)}(k) - P_{\pi_1(a)}(k))e(k) \in H_{\lambda}^{a}(\Gamma).$$

We obtain:

$$\Phi_{\lambda}^{a} \circ BK \circ JPA(g_a) = \sum_{i \in J_\lambda} \sum_{k' \in K_{j, a}^{\lambda}} \psi_{\pi_1}^{-1}(\psi_{\pi_1(a)}Q_{\pi_1(a)}(k) - P_{\pi_1(a)}(k))e(k)\psi_{\pi_1}^{-1} + \sum_{i \in J_\lambda} \sum_{k \in K_{\lambda}^{a}} \sqrt{q}\psi_{\pi_1}^{-1}e(k)\psi_{s_{a+1}}^{-1}.$$
We first focus on the first sum $S_1$: let $t \in J^n$ such that $t_a = t_{a+1}$ and $k \in K^t$. We can notice that thanks to Proposition 6.7 we have $\pi_t(a + 1) = \pi_t(a) + 1$. Using Lemma 6.20 and the properties of $\hat{Q}$ we have (recalling the notation $k^t$ introduced at (6.10)): 

$$Q_{\pi_t(a)}(k)e(k)\psi_{\pi_t} = \hat{Q}_{\pi_t(a)}(k)(y_{\pi_t(a)}, y_{\pi_t(a)+1})\psi_{\pi_t}e(k^t) = \psi_{\pi_t}\hat{Q}_{\pi_t(a)}(k)(y_a, y_{a+1})e(k^t) = \psi_{\pi_t}\hat{Q}_a(k^t)(y_a, y_{a+1})e(k^t).$$

The same proof gives $P_{\pi_t(a)}(k)e(k)\psi_{\pi_t} = \psi_{\pi_t}P_a(k)e(k^t)$. We have thus:

$$\psi_{\pi_t}^{-1}Q_{\pi_t(a)}(k) - P_{\pi_t(a)}(k))e(k)\psi_{\pi_t} = \psi_{\pi_t}^{-1}(\psi_{\pi_t(a)}Q_a(k^t) - P_a(k^t))e(k^t).$$

Using (1.19) and Lemma 6.22, we obtain:

$$\psi_{\pi_t}^{-1}(\psi_{\pi_t(a)}Q_a(k) - P_{\pi_t(a)}(k))e(k)\psi_{\pi_t} = \psi_{\pi_t}^{-1}(\psi_aQ_a(k^t) - P_a(k^t))e(k^t).$$

Finally, since $s_a \cdot t = t$, we get by Proposition 6.18:

$$\psi_{\pi_t}^{-1}(\psi_{\pi_t(a)}Q_a(k) - P_{\pi_t(a)}(k))e(k)\psi_{\pi_t} = (\psi_aQ_a(k^t) - P_a(k^t))e(k^t),$$

so the first sum becomes:

$$\sum_{t_a \in J^n} \sum_{k \in K^{t_a}} \psi_{\pi_t}^{-1}(\psi_{\pi_t(a)}Q_a(k) - P_{\pi_t(a)}(k))e(k)\psi_{\pi_t} = \sum_{t_a \in J^n} \sum_{k \in K^{t_a}} (\psi_aQ_a(k^t) - P_a(k^t))e(k^t).$$

We now focus on the second sum $S_2$: let $t \in J^n$ with $t_a \neq t_{a+1}$ and let $k \in K^{t_a}$, $k \in K^{t+a}$. By Lemma 6.21 we have directly:

$$\psi_{\pi_t}^{-1}(e(k)a_{\psi_{\pi_t}}) = e(k^t)\psi_{\pi_t}^{-1}a_{\psi_{\pi_t}} = e(k^t)\psi_{\pi_t}.$$

Thus:

$$\sum_{t_a \in J^n} \sum_{k \in K^{t_a}} \sqrt{Q}\psi_{\pi_t}^{-1}.e(k)a_{\psi_{\pi_t}} = \sum_{t_a \in J^n} \sum_{k \in K^{t_a}} \sqrt{Q}\psi_{\pi_t}^{-1}e(k^t)a_{\psi_{\pi_t}} = \sum_{t_a \in J^n} \sum_{k \in K^{t_a}} \sqrt{Q}\psi_{\pi_t}^{-1}e(k^t).$$

Finally, by (6.39)–(6.41) we get:

$$\Phi^\Delta \circ \overline{BK} \circ \overline{JPA}(g_a) = \sum_{t_a \in J^n} \sum_{k \in K^{t_a}} (\psi_aQ_a(k) - P_a(k))e(k) + \sum_{t_a \in J^n} \sum_{k \in K^{t_a}} (\psi_aQ_a(k) - P_a(k))e(k)

= \sum_{k \in K^n} \sum_{t_a \in J^n} (\psi_aQ_a(k) - P_a(k))e(k)

= \Phi^\Delta \circ \overline{BK} \circ \overline{JPA}(g_a) = \sum_{k \in K^n} (\psi_aQ_a(k) - P_a(k))e(k),$$

which is (6.38).

To conclude, we have checked that the algebra homomorphisms $\overline{BK}$ and $\Phi^\Delta \circ \overline{BK} \circ \overline{JPA}$ coincide on every generator of $\overline{\Gamma_{d,n}}(q)$, hence we have the following theorem.
Theorem 6.35. We have
\[ \bar{B}\mathcal{K} = \Phi^A \circ \bar{B}\mathcal{K} \circ J\mathcal{P}A, \]
and the diagram of Figure 3 commutes.

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