Review—Lithium Plating Detection Methods in Li-Ion Batteries

To cite this article before publication: Umamaheswari Janakiraman et al 2020 J. Electrochem. Soc. in press https://doi.org/10.1149/1945-7111/abd3b8

Manuscript version: Accepted Manuscript

Accepted Manuscript is "the version of the article accepted for publication including all changes made as a result of the peer review process, and which may also include the addition to the article by IOP Publishing of a header, an article ID, a cover sheet and/or an 'Accepted Manuscript' watermark, but excluding any other editing, typesetting or other changes made by IOP Publishing and/or its licensors."

This Accepted Manuscript is © 2020 The Author(s). Published by IOP Publishing Ltd.

As the Version of Record of this article is going to be/has been published on a gold open access basis under a CC 4.0 licence, this Accepted Manuscript is available for reuse under the applicable CC licence immediately.

Everyone is permitted to use all or part of the original content in this article, provided that they adhere to all the terms of the applicable licence referred to in the article – either https://creativecommons.org/licenses/by/4.0/ or https://creativecommons.org/licenses/by-nc-nd/4.0/

Although reasonable endeavours have been taken to obtain all necessary permissions from third parties to include their copyrighted content within this article, their full citation and copyright line may not be present in this Accepted Manuscript version. Before using any content from this article, please refer to the Version of Record on IOPscience once published for full citation and copyright details, as permissions may be required. All third party content is fully copyright protected and is not published on a gold open access basis under a CC licence, unless that is specifically stated in the figure caption in the Version of Record.

View the article online for updates and enhancements.
# Review—Lithium Plating Detection Methods in Li-Ion Batteries

| Journal:          | *Journal of The Electrochemical Society* |
|-------------------|-----------------------------------------|
| Manuscript ID     | JES-102473.R1                            |
| Manuscript Type:  | Review Paper                            |
| Date Submitted by the Author: | 15-Nov-2020                        |
| Complete List of Authors: | Janakiraman, Umamaheswari; General Motors Technical Center, Battery Cell Engineering Garrick, Taylor; General Motors Technical Center, Virtual Design, Development, and Validation; General Motors Corp, Virtual Design, Development, and Validation Fortier, Mary; General Motors Technical Center, Battery Cell Engineering |
| Keywords:         | Batteries, Batteries – Li-ion, Electrochemical Engineering |
Review—Lithium Plating Detection Methods in Li-Ion Batteries

U. Janakiraman,1,2 Taylor R. Garrick,2,* and M. E. Fortier1

1 Battery Cell Engineering, General Motors, Warren, Michigan 48093, USA
2 Global Virtual Powertrain & Electrification, General Motors, Warren, Michigan 48093, USA

*Electrochemical Society Member.
E-mail: umamaheswari.jp@gmail.com, uma.viswanathan@gm.com

Lithium-ion batteries (LiB) offer a low-cost, long cycle-life and high energy density solution to the automotive industry. There is a growing need of fast charging batteries for commercial application. However, under certain conditions of high currents and/or low temperatures, the chance for Li plating increases. If the anode surface potential falls below 0 V versus Li/Li\(^+\), the formation of metallic Li is thermodynamically feasible. Therefore, determination of accurate Li plating curve is crucial in estimating the boundary conditions for battery operation without compromising life and safety. There are various electrochemical and analytical methods that are employed in deducing the Li plating boundary of the Li-ion batteries. The present paper reviews the common test methods and analysis that are currently utilized in Li plating determination. Knowledge gaps are identified, and recommendations are made for the future development in the determination and verification of Li plating curve in terms of modeling and analysis.

Introduction

Li-ion batteries (LiBs) have been widely used in various applications like electric vehicles and energy storage systems, which require great durability and safety. For practical applications, fast charge capability of the battery is desired, which can be considered as 15 minute charging to 80% state of charge (SOC). However, certain charging conditions are known to provoke Li plating which can impact safety.\(^{[1-6]}\) Li plating is an undesired side reaction on the anode side of Li-ion batteries where Li ions are reduced to metallic Li instead of intercalating into the anode crystal structure. Deposition of metallic Li also enables dendrite formation. This phenomenon occurs when the local anode potential drops below 0 V vs. Li/Li\(^+\). In the presence of local inhomogeneity of the graphite anode surface, Li plating can be localized\(^{[7, 8]}\). It can also be induced by slow intercalation kinetics and high SEI (Solid Electrolyte Interface) film resistances in the anode, due to low temperatures or high current rates\(^{[9, 10]}\) as well as by slow diffusion processes\(^{[11, 12]}\). A poorly designed cell or choosing improper charging protocols may also induce Li plating. Anode material type has also been seen to influence Li plating due to the differences in equilibrium potential. For example, hard carbon is better than graphite and Li titanate is the least affected by Li plating\(^{[10, 13]}\). Similarly, electrolyte composition also affects the low-temperature performance of Li-ion batteries\(^{[9, 14-16]}\). In general, the deposited metallic Li is reversible. However, during oxidation, Li dendrites can lose contact with the active material, leading to dead Li in the cell\(^{[17]}\). Furthermore, the plated metallic Li is highly reactive and forms an irreversible SEI with the electrolyte\(^{[18, 19]}\). Both the dead Li as well
as the SEI can lead to a rapid capacity fade of the battery. Therefore, measurement of the cell capacity and resistance serve as the direct way to detect Li plating. Different studies suggest that metallic Li deposited on a graphite electrode is partly reversible. It has been reported that deposited Li can re-intercalate into the graphite during relaxation of the electrode.\textsuperscript{[11, 20]} Oxidation (‘stripping’) of metallic Li was observed during discharge processes. Both processes of Li recovery (re-intercalation as well as stripping) can be directly observed via the cell potential.\textsuperscript{[21-23]}

Methods for detection of Li plating can be divided into the following categories: (1) Measurement of anode potential vs Li/Li\textsuperscript{+} with a reference electrode.\textsuperscript{[24-27]} (2) Battery destructive physical analysis and imaging of anode.\textsuperscript{[28, 29]} (3) Electron Paramagnetic Resonance (EPR)\textsuperscript{[30, 31]} and Nuclear Magnetic Resonance (NMR)\textsuperscript{[32, 33]} to detect a particular range of resonance frequencies corresponding to metallic Li. (4) High precision measurement of coulombic efficiency\textsuperscript{[34, 35]} in which the plot of coulombic inefficiency per hour falls on the same curve for different tests if cell aging is dominated by SEI growth. Deviation from this curve indicates Li plating. (5) Measurement of cell thickness\textsuperscript{[36-39]} because Li metal deposition induces larger volume change than Li\textsuperscript{+} intercalation into graphite, estimated to be 0.37 cm\textsuperscript{3}/Ah\textsuperscript{[36]}. Thickness change becomes irreversible when Li plating occurs. (6) Measurement of graphite lithiation degree via in-situ Neutron Diffraction\textsuperscript{[40]}. (7) Detection of a distinct voltage plateau after charging\textsuperscript{[9, 21-23, 41-44]}. In addition to the above techniques, there are several new methods which have been developed in recent years such as (8) AC impedance analysis\textsuperscript{[45]} in which the plot of interfacial impedance versus SOC decreases with increasing SOC when Li plating occurs, (9) Accelerated Rate Calorimetry\textsuperscript{[46]} which measures the heat of reaction in the presence of plated Li, (10) Acoustic Time of Flight\textsuperscript{[47]} detection of transmitted ultrasonic waves as it changes with cell thickness during Li plating, and (11) Mass Spectrometry Titration\textsuperscript{[48]} to quantify the SEI components. The details of these different methods will be discussed in this review article.

The most common Li plating detection method is the detection of a voltage plateau due to the Li stripping process which indicates the occurrence of Li plating during charging. The voltage plateau can occur either at the beginning of discharge or during relaxation after charging. Among the above methods, the voltage plateau method is the most feasible method for online detection of Li plating in real-world applications, because it is nondestructive and does not require specialized and expensive equipment. Prior works demonstrated the feasibility of adopting the voltage plateau method as a method for Li plating detection. Therefore, mathematical models have been extensively applied to predict performance and life of LiBs and gain fundamental insights into the internal cell characteristics.\textsuperscript{[49-56]} Several works have been reported to study Li plating during overcharge\textsuperscript{[57, 58]} or during charge at high rates and/or low temperatures\textsuperscript{[33, 59]}, and to study the capacity loss induced by Li plating\textsuperscript{[60]}. Most of these models only considered the Li plating process and neglected the Li stripping process, and therefore cannot capture the voltage plateau behavior. Hein and Latz\textsuperscript{[61-64]} attempted to simulate Li stripping and the associated voltage plateau behavior using a 3D microscopic model. Through this work, it was shown that Li metal distribution is highly nonuniform in the anode.
In this review, we are offering an overview of the various Li plating detection and analysis methods, including mathematical models that are used for simulating the electrochemical characterization of the Li plating event.
Discussion

Three-Electrode Diagnostics

During charging at low temperatures, high rates, and high states of charge, the deposition of metallic Li on anodes occurs which leads to rapid battery aging and failure. This Li deposition on anodes can be detected in battery cells with a reference electrode. However, commercial cells in automotive or consumer electronics applications do not contain a reference electrode. Recently, much work has been done by constructing a representative small three-electrode setup (i.e. anode, cathode, and reference.)

Bhardwaj used a Li reference electrode built into a cell to investigate a charging strategy to avoid Li deposition on anodes. Zhang utilized a reference electrode while investigating a stepwise charging procedure with reduced current at the beginning of charging followed by a higher current. Itou collected electrodes from aged cells and reconstructed them together with a Li reference electrode and was able to determine the resistances of the anode and the cathode versus the reference separately. Waldmann constructed three-electrode cells from commercial 18650s with the addition of a Li reference electrode. This allowed for the measurement of the anode potential against the Li reference during charging of the reconstructed cells and determination of Li deposition. Waldmann observed from anode potential measurement that Li deposition on anodes in commercial cells is possible even at ambient temperature and low charge rates.

Koch has also done work on three electrode setups. These offer the ability to monitor the onset of Li plating in an electric vehicle and link with controller logic in a battery management system. However, to date, this type of reference electrode setup has not been seen in a vehicle application. Commercial cells in automotive or consumer electronics applications do not contain a reference electrode. Moreover, the location of a reference electrode should be carefully controlled to measure resistances at each electrode and solid-state diffusion rates. While the three-electrode method can determine when Li plating occurs, visual verification to quantify Li plating is still necessary.

Measurement of Coulombic Efficiency

When lithium plating occurs, a part of the deposited Li can intercalate into anode and remains in electrical contact with the negative electrode. Therefore, it can be stripped during the discharge step.

\[ Q_{stripping} = Q_{rev,plating} \]  

This reversible part does not cause any capacity loss because it is just another way of charging the graphite electrode. In contrast, a fraction of the plated Li can remain on the anode surface which may have fragile electrical contact with the negative electrode. Such Li is referred to as irreversible plated Li and may become electrically isolated from the graphite electrode during the stripping process. Irreversible plating, i.e. dead Li, is therefore
mainly responsible for the capacity losses during low temperature charging with high currents.\textsuperscript{[21]}

There are additional capacity losses due to the formation of surface films on the plated Li. Initially, Li is plated directly on the electrode surface.\textsuperscript{[75]} The plated Li can grow through the SEI and react with the electrolyte.\textsuperscript{[11]} Subsequent surface film formation causes capacity losses which are not distinguishable from irreversible plating. Moreover, the formation of such surface films can also cause electrical isolation of deposited Li.\textsuperscript{[59]} The charge corresponding to irreversible plating $Q_{\text{rev/plating}}$ includes the capacity loss caused by surface film formation. It is generally assumed that capacity losses due to SEI formation are negligible at low temperatures. The capacity loss during a charge/discharge cycle is simply the difference between the charge throughput during charge $Q_{ch}$ and discharge $Q_{dc}$, i.e.

$$Q_{\text{irrev/plating}} = Q_{ch} - Q_{dc}$$ \[2\]

The total charge during Li plating is

$$Q_{\text{plating}} = Q_{\text{rev/plating}} + Q_{\text{irrev/plating}}$$ \[3\]

The reversibility “R” of the plating process,

$$R = Q_{\text{rev/plating}}/Q_{\text{plating}}$$ \[4\]

The amount of Li plating,

$$n_{\text{plating}} = Q_{\text{plating}}/F$$ \[5\]

(Faraday constant $F = 96485.3399$ C mol$^{-1}$)\textsuperscript{[76]} as Li plating is a one-electron process ($\text{Li}^+ + e^- \rightarrow \text{Li}$). The mass of plated Li $m_{\text{plating}}$ is then calculated according to

$$m_{\text{plating}} = n_{\text{plating}} \times M_{\text{Li}}$$ \[6\]

(molar mass of Li, $M_{\text{Li}} = 6.94$ g mol$^{-1}$).\textsuperscript{[76]}

Reversibility “R” of the plating process can be determined by the quantitative detection method. The reversibility decreases towards high SOC. The following processes can take place at high SOC: formation of fragile Li deposit, subsequent chemical intercalation of some of the plated Li, and surface film formation on the residual plated Li. These processes lead to a poor electrical pathway and eventual isolation of fragile Li deposited during Li stripping. Such dead Li cannot be stripped and is therefore electrochemically inactive, i.e. the reversibility is reduced. Li plating increases at lower temperatures due to kinetic polarization effects. Therefore, more plated Li is prone to lose its electrical contact with the graphite electrode, which aggravates irreversible plating and reduces the reversibility. The coulombic efficiency, $CE$, describes the ratio between discharge and charge during one cycle, i.e.

$$CE = Q_{dc}/Q_{ch}$$ \[7\]
It typically illustrates capacity losses during charging, e.g. SEI formation at ambient temperatures. CE at low temperatures can be understood by the following equation,

\[ CE = \frac{Q_{ch} - Q_{irrev.plating}}{Q_{ch}} \]  

[8]

As the irreversible Li plating increases at low temperature, coulombic efficiency decreases. In a plot of SOC vs CE, a linear SOC dependence of the plated Li mass can be observed for low to medium SOC values. Decreasing temperature leads to higher Li plating growth rates and shifts the plating initiation point to lower SOC which means that the plating process starts earlier. It is crucial that the stripping plateau is clearly identifiable and not overlain by plateaus of the cathode material.

As described above, when Li plating occurs during cycling, the coulombic efficiency of the cell will decrease. Based on this premise, Li plating should be detectable as small changes to the coulombic efficiency of cells cycling at different rates. This implies that the detection limits for Li plating will be dictated by the precision and accuracy of measurements of the coulombic efficiency. A High Precision Charger\cite{34, 77} should therefore have a high degree of sensitivity for detecting Li plating in any cell type under standard cycling conditions. Burns et al carried out the High Precision Coulometry (HPC) study on a 3.4 Ah Panasonic 18650 (NCA/graphite) cells at different charge rates. They cycled the cells at 25°C and between 2.8 and 4.2 V at charge rates of C/10, C/5, C/3, C/2 or 1C and a discharge rate of C/10 on the High Precision Charger. A class “A” resistance temperature detector (RTD) (US Sensor, PPG101A1) was attached to the cells to measure variations in cell temperature at high rates. After cycling, cells were discharged to 0 V (to ensure safety while opening) and opened to look for evidence of Li plating on the negative electrode. The coulombic efficiency (CE) was time normalized to compare cells that are in the time (h) dependent regime. The coulombic inefficiency normalized over time (CIE/h) shows all cells at low rates fall on the same curve indicating, they are in a purely time dependent degradation regime which agrees well with the previous work from Smith et al.\cite{78, 79} However, when charging at 1C or higher the CIE/h becomes worse than those cells in the time dependent regime indicating the occurrence of Li plating. Based on the work of Smith et al.\cite{78, 79} it was shown that normally parasitic reactions are predominantly time dependent and cycle independent. The only exception is parasitic reactions on Silicon negative electrodes which are expected to be strongly cycle dependent due to the large volume changes during each cycle. Therefore, the increase in CIE/h must be attributed to Li plating as it is a cycle dependent degradation mode. As can be seen in Figure 1, with increasing rate, the amount of Li plating/stripping increases, and the CIE/h gets worse.
Figure 1. Capacity and Coulombic efficiency and inefficiency per hour (CIE/h) versus time for pouch cells cycling at 30°C during cycling with a two-step charge process from C/50 to 5C charge rates. Reprinted from J. C. Burns et al, Journal of The Electrochemical Society, 162 (6) A959-A964 (2015). [34] (CC BY, https://creativecommons.org/licenses/by/4.0).

From this data it can be seen when Li plating occurs it causes an increase in the degradation rate of the cells. The cells charged at a 5C rate decrease their capacity loss rate around 200 hours (near 170 mAh of remaining capacity) which is likely due to the large difference in over potential of the negative electrode relative to the positive electrode. When the cell is reaching the upper cutoff potential of 4.2 V, the state of charge of the graphite is significantly less than in the early cycles (hence the lower full cell capacity). As graphite is not being charged to as high an electrode state of charge, the voltage curve of the graphite no longer reaches 0 V vs Li/Li⁺ to initiate Li plating even with the same overpotential from the high current.

Figure 2 shows the results for CIE/h as a function of charge rate at different temperatures. Coulombic efficiency and CIE/h are taken as an average over the final ~50 hours of cycling while fade is taken as a linear fit to the capacity versus time data over that same final ~50 hours. Figure 2 shows that the CIE/h for cells cycled at lower rates at a given temperature is a constant. At lower temperature this constant is smaller as the degradation rate is less (than at higher temperature) due to the slower rate of parasitic reactions. However, as the rate is increased, the CIE/h begins to deviate from the low-rate constant value and becomes larger. This leads to “kinks” in the CIE/h versus rate curves where CIE/h is constant during the time dependent regime and then increases with rate during the charge rate dependent regime (where Li plating occurs). The “kinks” indicate the charge rate at which Li plating begins at that temperature. There are measurable differences in the fade rates (especially when the efficiency becomes very poor at high
rates) but the difference in fade rates is so small that it is much easier to resolve the onset of plating by examining the CIE/h. Burns et al further confirmed Li plating in their study, by opening the cells after cycling at different rates at ~50% state of charge and inspecting the anodes, in an argon-filled glove box.[34]

![Figure 2](image_url)

Figure 2. Charge C-rate vs CIE/h to 5C charge rates at different temperatures. Reprinted from J. C. Burns et al, *Journal of The Electrochemical Society*, 162 (6) A959-A964 (2015). [34] (CC BY, https://creativecommons.org/licenses/by/4.0)

Figure 2 shows a clear indication of the Li plating as opposed to the Li deposition/stripping curve from the dV/dt or dV/dQ plots. At 12°C, the “kink” in the CIE/h curve occurs between C/5 and C/2 which implies that the cell charged at C/5 should not show Li plating while the cell charged at C/2 should show small amounts and then the higher. Thus, high precision coulometry when coupled with a robust thermal chamber is an effective and non-destructive tool for Li plating determination.

Voltage Relaxation and dV/dt

In the plot of cell voltage vs time, the occurrence of voltage plateau is ascribed to stripping of deposited Li metal and thus indicates that Li is plated during charging. The voltage plateau can occur either at the beginning of discharge or of relaxation after charging. Among the various methods, the voltage plateau method is the most feasible method for online detection of Li plating in real-world applications, as it is nondestructive and does not require special and expensive equipment. Smart et al.[9,41] were among the first to apply discharge voltage plateau after low temperature (~40 °C) charging as a tool for Li plating detection and proposed that the length of plateau can be used as a measure of deposited Li metal amount. Petzl and Danzer[21] presented a differential voltage approach to detect and quantify Li plating. In their tests, a cell was charged with a 1C rate at < -20 °C and discharged immediately after charging. A voltage plateau appeared at the beginning of discharge. Later, Uehmann et al.[22] proposed to detect a voltage plateau during relaxation (rest time after charging) rather than during discharge, which is advantageous as the cell
has no net current at open-circuit, making the plateau more pronounced and easier to detect. In the work by Uhlmann et al. the Open Circuit Voltage (OCV) after charge showed plated Li can chemically intercalate into graphite when the cell is at rest. The process of Li plating followed by intercalation into graphite gives rise to a characteristic ‘dual plateau’ OCV feature. This plateau was interpreted as mixed potential resulting from the interplay of Li intercalation and dissolution of deposited metallic Li on the anode surface. The first potential plateau during relaxation corresponds to the mixed potential of the plated Li and lithiated graphite (Li$_x$C$_6$). As the plated Li chemically intercalates, the mixed Li/Li$_x$C$_6$ potential transitions to a pure Li$_x$C$_6$ potential after a threshold amount of surface Li has been depleted. This phenomena is more prominent as a voltage derivative (Figure 3), which was first proposed by Schindler et al. While differential OCV (dV/dt) have been used by many groups to study Li plating, the experimental study was either limited to cold temperatures (≤ 0°C) or charging to high graphite states of charge (SOC). However, EV battery control systems require a technique that indicates initiation of Li plating at high charge rates (C-rates) and at moderate temperatures typically used for fast charge.

![Figure 3: Typical voltage relaxation plot of a LiB as a function of % state of charge.](image)

The voltage relaxation method was adopted by many other researchers. Bauer et al. applied this method to study the impacts of charging temperature on Li plating and found that the duration of voltage plateau was longer at lower charging temperatures. This work is closer to real-world applications as the tested cell is a commercial large-size (16Ah) cell and at practical temperatures (-7.5 °C ~ +27.5 °C). Luders et al. further applied the voltage relaxation method along with in-situ neutron diffraction to examine the impacts of charge rate on Li plating, with a 1.95Ah 18650 cell charged at -2 °C. The increase of the Li plating amount during cell relaxation was estimated from the gain in graphite lithiation degree as measured by neutron diffraction and it was found to be proportional to the duration of voltage relaxation plateau. These results indicate that the time to reach dV/dt peak can be used for assessing the Li plating amount in different tests. The rate of Li stripping is affected by the capability of Li$^+$ intercalation into graphite.
and by the rate of Li\(^+\) diffusion and migration from the separator to the foil. Furthermore, the amount of Li plating is strongly influenced by the SOC and the measurement temperature.\(^{[81, 82]}\) It necessitates the verification of Li plating determination by voltage relaxation study with subsequent post-mortem analyses right after charging if reversible lithium plating is to be analyzed.\(^{[83]}\) The stripping plateau can be convoluted with other plateaus in a full cell configuration, including plateaus related to dissolution of overshooting graphite phases. This prompts the questioning of the reliability of the Li plating determined by the voltage relaxation analysis.\(^{[84]}\)

In addition to the voltage relaxation method (dV/dt), another widely reported technique for electrochemical Li detection is differential voltage analysis of the discharge profile (dV/dQ or dQ/dV). In the plot of derivative of voltage over capacity (dV/dQ), a distinct peak appears at the beginning of discharge. This peak indicates the end of Li metal stripping, and the discharge capacity at the dV/dQ peak corresponds to the total amount of Li plated during charging. This method can estimate the amount of plated Li based on the length of the plateau.\(^{[21, 23]}\) While dV/dt method requires a rest period after charge, dV/dQ must have immediate discharge to maximize the signal corresponding to Li stripping. The former is easier to implement whereas the latter can give false positive plating signals.\(^{[85]}\) Moreover, all the electrochemical methods discussed here have limitation to detect localized deposition due to averaging effects, which will be discussed later in this article.

**Li Plating Modeling**

Due to the detrimental effects that Li plating can have on battery performance, a push to develop mathematical methods to predict the onset of metallic Li growth has been seen in recent years. These numerical approaches through modeling offer a way to quantify the amount of plated Li which may offer a clearer picture into the extent of Li plating compared to the experimental in-situ and ex-situ methodologies, and offer the ability to be loaded into controllers\(^{[86, 87]}\) to provide an electrochemical model for in-vehicle applications. Arora and White may have been the first to propose a set of equations that describe the Li deposition reaction at overcharge.\(^{[57]}\) These equations were coupled with traditional porous electrode theory as detailed by Newman and colleagues.\(^{[88-95]}\)

Regarding Li transport in a Li ion battery, Li is assumed to transport from the solid matrix of one active material (such as the anode) governed by transport properties of the active material and reacts at the interface between the active material and the electrolyte governed by the electrochemical kinetics. Li salt in the electrolyte diffuses from one porous electrode, through a separator, to another porous electrode, governed by the liquid phase transport of the electrolyte. At the interface of the second electrode (such as the cathode) and the electrolyte, the Li reacts and intercalates into the solid phase. Traditional porous electrode theory governs these reaction and transport phenomena as is briefly detailed below.

It is assumed that Ohm’s law holds in the solid phase of a porous electrode

\[
i_1 = -\sigma_{\text{eff}} \frac{d\phi_1}{dx}
\]  \(^{[9]}\)
with boundary conditions such that charge is maintained and current flows from one porous electrode to another. When considering an electrochemical cell in one dimension, at the interface of one porous electrode and current collector:

\[
\sigma_1 \frac{d \phi_1}{dx} = i_{app} \tag{10}
\]

At the interface of the second porous electrode and current collector:

\[
\sigma_1 \frac{d \phi_1}{dx} = -i_{app} \tag{11}
\]

Charge is maintained in the solution phase of each porous electrode,

\[
i_2 = -\kappa_{eff} \frac{d \phi_2}{dx} + \left(\frac{2 \kappa_{eff} R T}{F}\right) \left(1 + \frac{d \ln f}{d \ln c_2}(1 - t_+)\right) \frac{d \ln c_2}{dx} \tag{12}
\]

as well as the separator,

\[
0 = -\kappa_{eff,s} \frac{d \phi_2}{dx} + \left(\frac{2 \kappa_{eff} R T}{F}\right) \left(1 + \frac{d \ln f}{d \ln c_2}(1 - t_+)\right) \frac{d \ln c_2}{dx} \tag{13}
\]

The material balance in the solution phase of a porous electrode is governed by Fick’s law:

\[
\frac{d c_2}{dt} = -\frac{d}{dx} \left(D_{eff} \frac{d c_2}{dx} + \frac{i_2 t_+}{F} \right) - s_i \frac{i_2 F}{n F} \frac{d i_2}{dx} \tag{14}
\]

As well as the separator:

\[
\frac{d c_{2,s}}{dt} = -\frac{d}{dx} \left(D_{eff,s} \frac{d c_{2,s}}{dx} \right) \tag{15}
\]

The material balance in the solid phase is calculated in a pseudo second dimension, hence why it is commonly referred to as a P2D model. Spherical particles are typically assumed and the solid phase Li concentration is modeled in 1D in the radial direction where:

\[
\frac{d c_1}{dt} = \nabla(D_1 \nabla c_1) \tag{16}
\]

\[
\frac{d c_{1,a}}{dr} \bigg|_{r=0} = 0 \tag{17}
\]

\[
-D_{1,a} \frac{d c_{1,a}}{dr} \bigg|_{r=r_p} = s_{i \alpha} \frac{r_p}{n F \alpha_2 \sigma_2} \tag{18}
\]

Finally, the electrochemical reactions are governed by the Butler-Volmer equation:

\[
i_2 = i_0 \left(\exp\left(\frac{a_\alpha F \eta}{RT}\right) - \exp\left(-\frac{-a_\alpha F \eta}{RT}\right)\right) \tag{19}
\]

Most treatment of Li plating using an electrochemical model use the set of equations listed above, with various modifications for the plating physics.
Arora and White proposed a set of equations that describe the Li deposition reaction at overcharge.\(^{[57]}\) The rate of the Li deposition reaction is charge-transfer-kinetic controlled and can be expressed by a Butler-Volmer expression as follows:

\[
i_k = i_{0,k} \left[ \exp \left( \frac{-\alpha_c k F}{RT} \eta_{s,k} \right) - \exp \left( \frac{-\alpha_a k F}{RT} \eta_{s,k} \right) \right]
\]

\[\text{[20]}\]

A number of different approximations can be made to simplify the computational processes while including the Li deposition side reaction. The cathodic Tafel expression can be used to describe the rate expression if either the deposition reaction is considered to be irreversible or if the amount of Li deposited is very small and reacts quickly with the solvent. In that case, the rate expression will be

\[
i_k = i_{0,k} \exp \left( \frac{-\alpha_c k F}{RT} \eta_{s,k} \right)
\]

\[\text{[21]}\]

Alternatively, the surface overpotentials may be sufficiently low so that the reaction can be expressed adequately using the linear approximation

\[
i_k = i_{0,k} \frac{(\alpha_a k + \alpha_c k) F}{RT} \eta_{s,k}
\]

\[\text{[22]}\]

Arora and White\(^{[57]}\) assumed that the Li deposition reaction is semi-reversible in that at least part of the deposited Li can dissolve during discharge. Some of the Li may react with the electrolyte to form insoluble product. In the above equations, \(i_{0,k}\) is the exchange current density and \(\eta_{s,k}\) is the local value of the surface overpotential defined as

\[
\eta_{s,k} = \phi_1 - \phi_2 - U_k - F j_{n,k} R_{film}
\]

\[\text{[23]}\]

where \(U_k\) is the open-circuit potential. The potential variables \(\phi_1\) and \(\phi_2\) represent the potentials in the solid and solution phases, respectively, and \(i_{0,k}, \alpha_a k,\) and \(\alpha_c k\) are the kinetic parameters. Based on the above equations and assumptions, the following kinetic expressions can be derived

\[
j_{n,1} = \frac{i_{0,1}}{F} \left[ \exp \left( \frac{\alpha_a k F}{RT} \left[ \phi_1 - \phi_2 - U_k - F j_{n,k} R_{film} \right] \right) - \exp \left( \frac{-\alpha_c k F}{RT} \left[ \phi_1 - \phi_2 - U_k - F j_{n,k} R_{film} \right] \right) \right]
\]

\[\text{[24]}\]

and

\[
j_{n,2} = \frac{i_{0,2}}{F} \left[ \exp \left( \frac{\alpha_a k F}{RT} \left[ \phi_1 - \phi_2 - U_2 - F j_{n,k} R_{film} \right] \right) - \exp \left( \frac{-\alpha_c k F}{RT} \left[ \phi_1 - \phi_2 - U_2 - F j_{n,k} R_{film} \right] \right) \right]
\]

\[\text{[25]}\]

where \(j_{n,1}\) and \(j_{n,2}\) correspond to the rates of Li intercalation and Li deposition reactions. The normal component of the current density is related to the pore wall flux by \(i_n = F j_n\). The open circuit potential \(U_2\) is equal to zero in this work because the potential is measured
with respect to a Li metal reference electrode in solution at the same local electrolyte concentration.

The interplay between the Li insertion reaction and the Li deposition reaction is governed by each set of kinetic parameters, resulting in the following equations governing the exchange current densities:

\[
i_{0,1} = F(k_{a,1})^{αc,1}(k_{c,1})^{αa,1}(c^0_1 - c_1)^{αa,1}(c_1)^{αc,1}(c_2)^{αa,1}\]

\[
i_{0,2} = F(k_{a,2})^{αc,2}(k_{c,2})^{αa,2}(c_2)^{αa,2}\]

where \(i_{0,1}\) is the exchange current density for the insertion reaction and \(i_{0,2}\) is the exchange current density for the Li deposition reaction. Considering a volumetric change due to Li deposition at the particle level, the exchange current density for the Li deposition reaction is directly related to the growth rate of a Li film on the particle surface:

\[
\frac{∂δ_{film}}{∂t} = -\frac{i_{n,2}M}{ρF}\]

where \(δ_{film}\) is the film thickness composed of solid Li and other products and \(M\) and \(ρ\) are the molecular weight and density of Li and products.

Tang\textsuperscript{[96]} observed that Li deposition occurs preferentially at the edges of the electrodes and that extending the negative electrode past the edge of the positive electrode can hamper this method of Li plating. A simplified mathematical model was used to simulate this behavior and probe the conditions under which this deposition occurs. This work is a critical extension of the traditional porous electrode model and considered edge effects in the electrodes. Through this work, it was determined that the extension of the negative electrode beyond the edge of the positive electrode provides the excess capacity at the critical location and prevents the Li plating from occurring. Tang\textsuperscript{[96]} assumed small currents and linearized the Buler-Volmer expression:

\[
i_2 = i_0 \frac{F}{RT} η\]

and considered the exchange current density, \(i_0\), as a function of concentration of Li in the solid phase and electrolyte:

\[
i_0 = Fk(c^0_1 - c_1)^{αa}(c_1)^{αc}(c_2)^{αa}\]

Although a galvanostatic charge means that the total current applied to the electrode is constant with time, the current density along the electrode is non-uniform. This means that a constraint of the system is an integral boundary and is defined as

\[
Φ_{electrode} i_n ds = l_{set} l_+\]

where \(l_+\) is the length of the positive electrode and \(l_{set}\) is the set length.
where \( l_p \) is the length of the positive electrode, and \( I_{set} \) is the average applied current density in amperes per meter squared, determined by dividing the theoretical capacity of the positive electrode by the time of charge

\[
I_{set} = F \frac{C_{p} + \rho + w + \Delta y}{t_{charge}} \tag{32}
\]

Other groups\[^{33, 97, 98}\] have proposed using two different Butler-Volmer equations to build on the work of Arora\[^{57}\] and White, many of which assume that the open circuit voltage of the Li plating reaction is zero, even though it is known that Li deposits prior to 0V vs. a Li reference electrode\[^{99-107}\].

Ren\[^{97}\] and colleagues proposed separate kinetics equations for Li plating and Li stripping, with the Li stripping reaction having the same kinetics as the Li plating reaction. When the overpotential is less than 0, Li plating will occur governed by Eq. 33

\[
j_{n,2} = \frac{i_{0,2}}{F} \left[ \exp \left( \frac{\alpha a_{2} F \eta_{Li}}{RT} \right) - \exp \left( - \frac{\alpha c_{2} F \eta_{Li}}{RT} \right) \right] \tag{33}
\]

When the overpotential is greater than 0 and some quantity of Li that can be removed is present, Li stripping will occur governed by Eq. 34

\[
j_{n,3} = \frac{i_{0,2}}{F} \left[ \exp \left( \frac{\alpha a_{2} F \eta_{Li}}{RT} \right) - \exp \left( - \frac{\alpha c_{2} F \eta_{Li}}{RT} \right) \right] \frac{\beta n_{Li,rev}}{1 + \beta n_{Li,rev}} \tag{34}
\]

In Eq. M26, \( \beta n_{Li,rev} \) is a correction term limiting the reaction to the amount of reversible Li. \( \beta \) is a large value, set to 1000 in Ref. \[^{97}\], such that the correction term is equal to 1 when a significant amount of plated Li is present, and 0 when \( n_{Li,rev} = 0 \). In both Eq. 33 and Eq. 34

\[
i_{0,2} = F k_2 (c_2)^{\alpha a_{2}} \tag{35}
\]

and \( \eta_{Li} \) is consistent with Eq. 23.

More recently, Baker and Verbrugge have developed a Multi-Species, Multi-Reaction (MSMR) Model\[^{108-112}\] for porous intercalation electrodes that has been applied to plating and dissolution of Li in a Li ion cell\[^{112}\]. In the latter work, a set of equations was derived to describe the overcharge of a Li ion cell. Two critical additions to published models were added to address Li plating. First, the handling of Li deposition and dissolution electrochemical reactions was addressed. Secondly, the plated Li in system can directly interact with the vacant sites in the anode, which will impact the system response. By including these theoretical developments with the MSMR model, the authors were able to illustrate Li intercalation and plating within a single particle model. One critical observation regarding the thermodynamics of the Li plating reactions shows that the open circuit voltage of the plating reaction with respect to a Li reference cannot be zero, as commonly assumed. Rather, it is only zero when the surface coverage of plated Li is large enough for the Li to interact with the electrolyte in the same manner as a bulk metallic Li electrode. Consequentially, when small surface coverages of Li on the surface of the anode particles exists, the plated Li interacts with the anode active material, resulting the OCV of
the plated Li being positive, which the authors detail in Eq. 11 in Ref. [112]. This provides an update to the historical assumption of Li deposition at 0V vs a Li reference.

In the Multi-Species, Multi-Reaction (MSMR) Model[108-112] a few parameters can be associated with the physical chemistry of the intercalation material enabling a quantitative fit of the open-circuit potential. In this model, different galleries of intercalated Li are assumed to be in thermodynamic equilibrium, so that their open circuit potentials are equal to some common equilibrium potential of the intercalation material. For each gallery $j$

$$U_1 = U_j^0 + RT \frac{\omega_j}{F} \ln \left( \frac{x_j - X_j}{x_j} \right)$$  \[36\]

$$x_j = \frac{x_j}{1 + \exp \left[ \frac{F(U_j - U_j^0)}{RT \omega_j} \right]}, \quad x(U_1) = \sum x_j$$  \[37\]

$$\frac{dx}{dU_1} = -\sum_{j=1}^{n} \frac{F}{RT \omega_j} \left[ x_j \exp \left( \frac{F(U_j - U_j^0)}{RT \omega_j} \right) \right]$$  \[38\]

$$i_{s,j} = i_{0,j} \left[ e^{\frac{(1-\beta_j)F}{RT} \eta_s} - e^{\frac{-\beta_j F}{RT} \eta_s} \right], \quad j_1 = \sum i_{s,j}$$  \[39\]

$$\eta_s = \phi_1 - \phi_2 - U_1(x)$$  \[40\]

$$i_{0,j} = i_{0,j}^{\text{ref}} \frac{\omega_j}{\beta_j} (X_j - x_j)^{\omega_j(1-\beta_j)} \left( \frac{c_2}{c_0} \right)^{1-\beta_j}$$  \[41\]

where $U_1$ is the open circuit voltage of the intercalation reaction, $\omega_j$ is a thermodynamic parameter fit to the open circuit voltage, $X_j$ is the maximum mole fraction of species $j$, $x_j$ is the mole fraction of Li in the solid phase, and $\beta_j$ is the symmetry factor for the $j$th intercalation reaction. The kinetics for the intercalation reaction are then governed by Eq. 39, a modification of Eq. 20 using the symmetry factor $\beta_j$.

The kinetics of the Li plating reaction are then given by

$$j_2 = i_{0,\text{Li}} \left[ e^{\frac{(1-\beta_{\text{Li}})F}{RT} \eta_s} - e^{\frac{-\beta_{\text{Li}} F}{RT} \eta_s} \right]$$  \[42\]

where

$$i_{0,\text{Li}} = i_{0,\text{Li}}^{\text{ref}} a_{\text{Li}}^{\beta_{\text{Li}}} \left( \frac{c_2}{c_0} \right)^{1-\beta_{\text{Li}}}$$  \[43\]

where $a_{\text{Li}}$ is the activity coefficient of Li and is related to the Li surface concentration and is defined as

$$a_{\text{Li}} = 1 - e^{-\Theta}$$  \[44\]
where \( \Theta \) corresponds to the ratio of the concentration of surface Li compared to a reference value. Combining Eq. 42 – 44 and replacing \( \Theta \) with \( U_2 \), the open circuit voltage of the plating reaction yields a kinetics expression that is a function of potentials, electrolyte concentration, and thermodynamic properties:

\[
j_2 = i_{0,Li}^{ref} \left( \frac{c_2}{c_2^0} \right)^{1-\beta_{Li}} e^{-\frac{F}{R}U_2} e^{\frac{F}{R}(1-\beta_{Li})(\phi_1-\phi_2)} - e^{\frac{F}{R}(-\beta_{Li})(\phi_1-\phi_2)} \]

This treatment removes the need for a correction term as seen in Eq. 44 which must be arbitrarily set.

Electrochemical modeling techniques are all based on a set of highly non-linear equations that describe the transport phenomena and electrochemical physics within the battery. While Thevenin\cite{113-127} models or electrochemical impedance spectroscopy models\cite{128-131} (EIS) are popular for fast predictions for electrical performance, they do not enable the user to couple the Li plating physics necessary for analysis. Retaining aspects of the porous electrode model is necessary for Li plating simulation. In order to speed up solution times, many different approximations have been made to simplify the equations and associated solutions depending on the accuracy and speed necessary in a given application.

The prevailing approach for increasing speed for electrochemical models without sacrificing accuracy is to use various mathematical techniques to speed up the solution time for the partial differential equations, or to convert the partial differential equations into a form that can be integrated faster through the use of numerical solvers. Subramanian\cite{132-136} developed various reduction methods that enabled real-time simulations. White has applied various methods such as orthogonal decomposition\cite{137}, state-space reformulation\cite{138, 139}, reducing the system of partial differential equations to ordinary differential equations\cite{139, 140}, and applying polynomial representations of the Li ion concentration in the solid phase\cite{141}. In recent years, the popularity of these models\cite{142-153} has increased due to their relative speed and accuracy.

Due to the development of fast and accurate reduced order models, the potential exists for real time Li plating estimation within a micro-controller. Zhao\cite{154} utilized a reduced order model\cite{155-157} to consider Li plating evolution. In this work, the ion concentration in the electrode was reduced using a polynomial approach, the ion concentration in the electrolyte was reduced using the state-space approach, and the electrochemical kinetics (i.e. Butler Volmer Equation) was linearized. Similar to previous approaches, the overall electrochemical reaction was split between the Li intercalation and Li deposition reactions

\[
j_{\text{total}}^{Li} = j_{\text{Li}}^{i} + j_{\text{Li}}^{\text{plating}} \tag{46}
\]

\[
j_{\text{Li}}^{\text{plating}} = a_s i_{0,p} \left[ \exp \left( \frac{\alpha_{a,p} F \eta_p}{RT} \right) - \exp \left( - \frac{\alpha_{c,p} F \eta_{Li}}{RT} \right) \right] \tag{47}
\]

where \( i_{0,p} \) is the exchange current density of the Li deposition reaction and was fit as a function of temperature, \( \alpha_{a,p} \) and \( \alpha_{c,p} \) are the dimensionless anodic and cathodic charge...
transfer coefficients and were assumed to be 0.33 and 0.67 respectively, due to the assumption of semi-reversibility with the Li deposition reaction. $\eta_p$ was calculated similarly to Eq. 23. The authors calculated the Li plating reaction rate (Figure 3d in Ref. [154]) and illustrate the increase in plating rate at lower temperatures, consistent with experimental studies. Choe has applied this reduced order model to investigate degradation with the intent of mitigating aging and Li plating due to fast charging.\textsuperscript{[144, 154-156, 163-167]} This type of approach could be utilized in a battery management system or charge controller to avoid Li plating completely, or to allow for a quantifiable amount of Li plating that could be electrochemically removed later.

Various research groups have approached integration of a full order model or reduced order model into a micro-controller. Perkins\textsuperscript{[58]} developed a control-oriented cell degradation model by deriving a reduced-order model for the Li plating deposition reaction occurring on overcharge. This work illustrated the differences in accuracy between a full-order model and the reduced order model, while allowing for a 5000x increase in solver speed, illustrating the theoretical capability of loading a ROM onto a micro-controller.

Florentino\textsuperscript{[168]} extended Perkin’s work and explored the use of controllers for managing battery cell performance using reduced order electrochemical models. In this work, model predictive control was utilized as a smart sensor and applied to a balancing architecture. As a smart sensor, the model predictive controller is transparent during the operation of a cell, provided that the cell does not enter into a state that would result in aging or Li plating. Once a condition is detected that would result in the cell being prone to Li plating, the model predictive controller intervenes and ensures that the cell does not enter said state.

Subramanian\textsuperscript{[133, 135, 136, 169-175]} has spent decades focused on modeling methods, and has more recently focused on reformulating\textsuperscript{[133, 135, 136, 173]} models for increased computational efficiency, developing battery management systems\textsuperscript{[172, 175]} using physics based reformulated models, and applying nonlinear model predictive control\textsuperscript{[174, 175]} to Li ion batteries. Subramanian’s work has illustrated the tradeoff between computational complexity and accuracy, the need for reduced order models in production in both the automotive and grid storage space, and has provided ample proof that physics based models are the next step in bridging the gap between scientific understanding of aging phenomena and the tools that are used for controls every day.

To detect Li plating, the first adopted approach was based on two-electrode models\textsuperscript{[176]} where the Li-ions transport rates in and out of the electrodes were used as indicators of the same. However, the approach doesn’t estimate the accurate Li-ion concentration. This limitation was overcome by using the extended model and cascaded observers. Dey et al\textsuperscript{[175]} expressed the insertion and extraction rates of the Li-ions as

$$E_p = \left| \frac{\partial CP_{avg}}{\partial t} \right|, I_N = \left| \frac{\partial CN_{avg}}{\partial t} \right|$$ \[48\]

where $E_p$ and $I_N$ are the Li-ion extraction rate from positive electrode and insertion rate into negative electrode. $CP_{avg}$ and $CN_{avg}$ are the volume-averaged Li-ion concentrations in the positive and the negative electrodes, respectively, and given as...
\[ CP_{avg} = \frac{1}{3\pi R^3} \int_{0}^{R^+} 4\pi r^2 C^+_s(r,t) dr \]  \[ CN_{avg} = \frac{1}{3\pi (R)^{-3}} \int_{0}^{R^-} 4\pi r^2 C^-_s(r,t) dr \]  

Li plating is detected when the extraction rate \( E_P \) exceeds the maximum possible extraction rate \( E_{P,max} \) and \( I_N \) drops below the minimum possible insertion rate \( I_{N,min} \). Their recommend ways to collect \( E_P \) and \( I_N \) data are: (1) Monte Carlo simulation studies on the battery model on varying different operating conditions, different levels of measurement noise, and different levels of modeling uncertainties or (2) experimental studies on a physical battery under different operating conditions[8, 43]. Then, choose the maximum and minimum of \( E_P \) and \( I_N \) as the limits. State and Capacity Estimation were performed for constant current and dynamic discharge conditions. The capacity fade is simulated by adding a side reaction current component into the plant-model.

The simulation scenario used for this case study is C/10 charging at ambient temperature, 20°C. To illustrate the Li plating condition in simulation scenario a step like additive C/5 charging input current component was used to the positive electrode of the plant-model. This additive input current causes a sudden change in current density in the positive electrode, while the current density in the negative electrode is kept constant. This essentially represents a loss in active Li-ions, as the Li-ion extraction rate from the positive electrode would be higher than the insertion rate in the negative electrode. In this case study, the positive electrode current density increases at \( t = 600 \) s. Using this approach, the bulk SOC could be tracked even after the occurrence of the plating.

Dey et al also studied the effect of measurement noise via simulation by using different noise levels in the all three measurements (voltage, temperature, and current). It is found that the scheme generates reasonable estimates approximately up to 25mV voltage measurement noise and 50mA current measurement noise. Above those ranges, the steady-state estimation error trends significantly higher reducing the effectiveness of the scheme. Other than measurement noise, measurement bias was also considered. For voltage and current measurements, the scheme performs satisfactorily up to 5mV and 12mA measurement bias. Beyond these values, the performance of the scheme degrades. Above 0.5 °C noise level in the temperature measurement, the estimates tend to diverge after a certain point when the signal-to-noise ratio gets smaller. The simulation results are yet to be correlated with the experimental measurement.

**Electrochemical Impedance Spectroscopy**

Multiple electrical and chemical processes occur during operation of a Li-ion battery including Li-ion diffusion in the electrolyte, migration through the SEI layer, charge transfer through the electrode/electrolyte interface, solid-state diffusion in the bulk of active material and electron transfer external to the battery via the current collectors[39, 40, 178]. These processes introduce impedance to the charge flow and cause a voltage drop between the two electrodes. Electrochemical Impedance Spectroscopy (EIS) is an efficient tool to identify these processes because each process has a different time constant or frequency range of excitation. [36, 179]
The decrease in EIS was first correlated with Li plating by S. Shindler et al. In their study they used a commercial lithium iron phosphate (LFP)/graphite cells. Voltage relaxation coupled with EIS was employed to detect lithium plating. Two main features were observed in the EIS namely a decrease of the high frequency intersection resistance and a respective decrease in the diameter of the semicircle representing the anodic charge transfer process. Both features have shown to be most pronounced for cases where the occurrence of lithium plating has been proven by the voltage relaxation method. In a recent publication, Koleti et al. studied the impedance analysis on a commercially available 18650-type Li-ion cells with a rated capacity of 3.1Ah. EIS is shown in Figure 4a. In figure 4b, the impedance due to solid diffusion limitations are referred to as $Z_{SDL,NE}$ or $Z_{SDL,PE}$, and the corresponding electrode surface potential is represented by $V_{NE,S}$ or $V_{PE,S}$; NE and PE stands for negative and positive electrodes, respectively. Figure 4c shows the EIS plots of the same cell at different SOC levels. As it can be observed, during charge/discharge, impedance in the mid-frequency range varies with SOC. Both $Z_{CTL}$ and $Z_{SEI}$ contribute to this range. As no significant change in the SEI layer thickness is expected in a single charge/discharge cycle the mid-frequency impedance variations can be attributed largely to the $Z_{CTL}$. The electrochemical models that describes the CTL through the Butler-Volmer equation also support this rationale. At low and high and high lithiation levels (0% and 100%), $Z_{CTL}$ is high and reduces as the lithiation level moves towards 50%. Therefore, the impedance changes in the mid-frequency region comes mainly from the $Z_{CTL}$ and tracking the battery impedance should indicate the way $Z_{CTL}$ is changing while the battery is being charged or discharged. The impedance measured at the transition frequency (referred as ftr which is marked with “*” in Figure 4c) at which impedance changes to 45° slope allows quantifying the total impedance from both the high and the mid-frequency regions. Since the impedance from the high-frequency region is nearly constant (as seen from Figure 4c), tracking the transition frequency impedance (referred as ZTR) allows tracking the impedance from the charge transfer limitations, $Z_{CTL}$.

$$Z_{TR} = Z_{\Omega} + Z_{SEI} + Z_{EL}$$  \[51\]

Where $Z_{\Omega} = Z_{\Omega1} + Z_{\Omega2}$ and $Z_{CTL} = Z_{CTL,NE} + Z_{CTL,PE}$

Here, $Z_{\Omega1}$ and $Z_{\Omega2}$ are Ohmic resistances at the NE and PE respectively; $Z_{CTL,NE}$ and $Z_{CTL,PE}$ are impedance due to charge transfer limitations at negative and positive electrodes, respectively. Figure 4d presents the ZTR profile against the battery SOC where ZTR drops initially as battery SOC raises from 0% and reaches a near flat level in the mid-SOC range before rising again as SOC goes further towards 100%.
Figure 4: Battery Impedance (a) EIS plot at 50% SOC (b) Equivalent electrical circuit diagram representation of the battery impedance (c) EIS plots at different SOC levels and (d) Impedance and resistance at the transition frequency, \( f_{tr} \). Reprinted from U.R. Koleti et al., *Journal of Power Sources* 451, 227798 (2020). Copyright (2020) with permission from Elsevier.

The method to track the impedance while charging/discharging is similar to the hybrid pulse power characterization (HPPC) test often used in developing equivalent circuit models\(^{[185-188]}\). Charge current is interrupted every 1% SOC increase for a predefined time period. While charging at low C-rates or during high ambient temperatures, differences between the lithiation levels of the electrode surface and the bulk are known to reduce\(^{[21-23, 159, 183]}\) due to the improved diffusion in the electrodes. Therefore, charging conditions that allow lithiation levels closer to their average can produce a ZTR profile similar to Figure 4d where the lithiation differences in the electrode particles are deemed to be negligible. Conversely, at high C-rates of charging or at low temperatures, the lithiation increases on the negative electrode because of reduced solid diffusion. Therefore, at high C-rates impedance decreases at a lower SOC level compared to charging at low C-rates or high temperatures. With further charging, lithiation tends to increase towards a full lithiation level, as the negative electrode voltage comes closer to Li reference potential well before the battery reaches 100% SOC. Additionally, the potential drop due to charge transfer limitations increases with increasing C-rate or decreasing temperature. The negative electrode potential can fall below the Li reference where Li plating commences. Once Li plating begins, electrode lithiation increases and \( Z_{CTL} \) decreases because the charging current is now divided into intercalating current and Li plating current. Additionally, as discussed within Ref. \(^{[97]}\) and \(^{[33]}\), the Li plating reaction at the NE also follows the Butler-Volmer equation that introduces a new branch of impedance in parallel to the existing \( Z_{CTL,NE} \). Similar to intercalating current that faces \( Z_{CTL,NE} \), Li plating current faces an impedance referred here as \( Z_{CTL,Li} \). As a result, the overall impedance at the NE drops after the onset of Li plating. On this basis, tracking the ZTR will allow detecting the onset of Li plating during the charge event. In the case of a charging event without Li
plating, ZTR follows the curve from Figure 4d where it faces continuously decreasing reduction at lower SOC levels to reach a near flat level in the mid-SOC range and then accelerates as SOC rises towards 100%. In the case of a charging event with Li plating, ZTR faces a second downtrend after the onset of Li plating because of the introduction of parallel impedance ($Z_{\text{CTL,li}}$) across the existing $Z_{\text{CTL,NE}}$. Detecting this second downward trend by tracking the ZTR shall, therefore, indicate the onset of Li plating. Figure 5 shows ZTR as a function of rate for test temperature of 20°C.

![Figure 5](image-url)

**Figure 5.** ZTR profiles without thermal cooling (a) ZTR profiles at different C-rates and (b) temperature profiles measured on the cell surface. Reprinted from U.R. Koleti *et al.,* Journal of Power Sources *451*, 227798 (2020). Copyright (2020) with permission from Elsevier.

In practical applications, temperature rise during a charging event is inevitable which can decrease the impedances due to electrolyte, CTL and SEI layer. As Li plating typically occurs towards the end of CC charging where temperature changes are not assumed to be significant, temperature influence on the internal processes may be
Therefore, Li-plating detection using the ZTR profiles is possible even under environmental conditions where temperature is increased during the charge process.

**Cell Thickness Measurement**

Bitzer and Grule\textsuperscript{[36]} successfully employed cell thickness measurement to detect Li plating. It is based on the concept of anode (and cell) volume expansion. During the LiB charge/discharge process, the active materials vary in volume, depending on the degree of Li-ion intercalation. Graphite anodes, for example, gain about 10% in volume while changing from the de-lithiated elementary state to fully lithiated LiC\textsubscript{6}\textsuperscript{[189]}. On the other hand, the cathode materials have a much smaller dilatation coefficient, for example, it is about 1% for NMC\textsuperscript{[190]}. Therefore, generally the thickness of LiB is governed by the anode volume change. One exception is Li titanate anode (LTO), which has a negligible volume change of only 0.1% - 0.3\%\textsuperscript{[190, 191]}. During charging, Li occupies the interstitials into the graphite lattice. Hence, the volume of the intercalation compound is smaller than the volume of the two separate materials. Li plating on the anode surface instead of intercalating into the anode leads to an additional gain in the overall volume. On the contrary, the de-intercalation from the cathode has no significant effect on the cell thickness. For 1 Ah of charge from the cathode to the anode, the amount of Li transferred is:

\[
n_{Li}(1Ah) = \frac{C}{eN_A} = 37.31 \text{ mmol} \tag{52}
\]

\*n_{Li}*: amount of Li substance [mol], \*N_{A}*: Avogadro constant [1/mol], \*e*: elementary charge [As]; \*C*: transferred charge [As]. When the anode is fully charged, every six carbon atoms of graphite takes one Li atom. Thus for 1 Ah of charge an amount of 6 nLi = 223.9 × 10^{-3} mol of carbon is consumed. With the molar volume of carbon in hexagonal structure of 5.31×10^{-6} m\textsuperscript{3} mol\textsuperscript{-1} its volume is 1.189 cm\textsuperscript{3}. According to the literature, carbon-based anode materials expand about 10% due to the intercalation of Li\textsuperscript{[189]}. Thereby charging 1 Ah causes a volume change of:

\[
\Delta V_C(1Ah) = \alpha_{LiC_6} n_{Li}(1Ah) 6 V_{m,C} = 0.12 \text{ cm}^3 \tag{53}
\]

\*\Delta V_C*: Change of the carbon volume due to Li intercalation [m\textsuperscript{3}]; \*\alpha_{LiC_6}* coefficient of expansion changing from C to LiC\textsubscript{6} =10\%; \*V_{m,C}*: molar volume of carbon [m\textsuperscript{3} mol\textsuperscript{-1}]. Given the molar volume of Li is, \*V_{m,LI}*: 13.02 × 10^{-6} m\textsuperscript{3} mol\textsuperscript{-1} the volume of Li corresponding to 1 Ah of charge is 0.49 cm\textsuperscript{3}. In the case of Li metal disposition on anode surface, instead of intercalation, the resulting expansion equals the volume of the Li in metallic form minus the volume the anode would gain due to the intercalation, as shown in Equation (2).

\[
V_{exp}(C_{pl}) = \frac{C_{pl}}{eN_A} (V_{m,LI} - \alpha_{LiC_6} 6 V_{m,C}) \tag{54}
\]

\*V_{m,LI}*: molar volume of Li [m\textsuperscript{3} mol\textsuperscript{-1}]; \*V_{m,C}*: molar volume of carbon [m\textsuperscript{3} mol\textsuperscript{-1}]; \*C_{pl}*: amount of plated Li. According to Equation 54, an amount of 1 Ah of plated Li causes a gain in volume of \*V_{exp}*(1 Ah) = 0.37 cm\textsuperscript{3}. Using this equation, it is possible to calculate the thickness of plated Li, knowing the surface area of electrode in the LiB. Thus, accurate measurement of cell thickness can serve to identify the onset of Li plating. Change of cell
thickness during standard cycling showed that the volume change due to Li intercalation into graphite takes place in three steps with different slopes which is highly reproducible. When the battery is charged with high current to cause Li plating, cell thickness increase could be observed. If the cell is at rest after plating, part of the plated Li intercalates into the graphite, causing some decrease in thickness. However, at higher C rates, the plated Li might not have sufficient time to re-intercalate to graphite, causing a thickness growth.

There are some practical challenges that affect the measurement accuracy. Due to cell inhomogeneity, dendritic deposition and gas evolution, the actual measured value is expected to be larger. Charging the cell with continuous currents leads to shifts in the Δcell thickness vs Ah curve. Similarly, measurement temperature can interfere with accuracy. However, at a temperature increase of 1 K at the cell surface the thermal expansion is smaller than the resolution of the dial thickness indicator. In addition to the temperature effect, the measurement point of thickness causes also affects measurement accuracy. Due to an existing inhomogeneity within every cell, the position of the measurement point plays a crucial role. Many pouch cells are known to start plating in a narrow area near the cell’s margin. When a wrong measurement point is chosen the experiment could potentially fail to detect the onset of Li plating. The best point is always cell specific and must be determined prior to the thickness measurements, e.g., by opening the cell. Pouch cells especially have problems arising from gas evolution originating from side reactions which interfere with accurate thickness measurement. Applying pressure on the cell by using appropriate spring can mitigate the problem and reduce the error. It was shown that during standard cycling other influences like the cathode expansion, temperature or gas evolution are small.

Optical Detection Methods

The most common method to detect Li plating is to open the cell after cycling and look for Li deposit. Generally, the anode from the plated cell is inspected by one or more of the following techniques such as optical microscope, SEM, TEM, AFM, XPS, Raman etc. One challenge with these ex situ techniques is, the plated Li can also intercalate back into the graphite which sometimes makes it challenging to visually observe microscopic Li plating. Moreover, this method is conducted ex-situ which can introduce sample preparation error. Nevertheless, optical detection methods are effective tools to identify Li plating, its morphology and composition.

SEM investigations\textsuperscript{192} enable the analysis of plated Li and dendrites by both in situ and ex situ. Gireaud \textit{et al.}\textsuperscript{193} observed dendrite growth on a Li surface at 20°C in Li/Li-coin cells ex situ. Their work showed the influence of the surface structure on the dendrite growth. In contrast, Sagane \textit{et al.}\textsuperscript{192} analyzed dendrite growth on copper under vacuum inside an in-situ SEM cell with a Li/LiPON/LATP/LiPON/Cu-configuration. They correlated the dendrite size with the applied current density. High current densities led to the formation of many small dendrites; low current densities led to the formation of fewer but larger dendrites. Steiger \textit{et al} studied the growth of mossy Li by in situ light microscope and ex situ SEM.\textsuperscript{194} They found that the mossy Li forms at the grain boundaries by the insertion of Li into defect sites. J. Guo \textit{et al} studied the effect of VC-LiNO$_3$ additive on Li metal by EIS and ex situ SEM and XPS.\textsuperscript{195} They constructed a cell with Li and copper
and used the electrolyte 1M LiPF$_6$ EC/DMC (1/1) as their baseline and studied the effect of additives, 2 vol.% VC and 0.1M LiNO$_3$. It was found that the selected additive combination contributes to positive modifications of SEI film on the surface of Li anode in carbonate electrolyte system. Nearly 100% efficiency was observed for Li plating/stripping in the presence of these additives and a very homogeneous morphology of cycled Li were obtained.

Harris et al used in situ optical microscope detection of Li plating on a full operational cell and observed color change of graphite anode as a function of Lithiation. Uhlman et al also used an optical test cell-housing to observe the graphite anode in situ via confocal microscopy whilst applying current. They used a charging current of 10 C for the test and. They observed the color change of graphite with lithiation.

Depending on the SOC, the graphite changes its color from black, red, to gold. This corresponds to the change of the band structure, due to the variation of the lattice spacing. Though this change happened very quickly in this measurement, due to the high charging current, it could be seen within the optical test-cell. After turning golden, which indicates the highest Li intercalation stage of graphite, a gray coating starts to cover the electrode. At the beginning of this process some small dots appear on the particles. Within 1-2 s, the whole observation area is covered by a gray structure which is interpreted as a layer of metallic Li. When observing the frontier between plated and non-plated surface area during the relaxation phase, several facts were recognized. On one hand, the plated area decreases while the frontier of the two areas moves back. This is due to dissolution of plated Li. With decreasing Li surface concentration, intercalation becomes favorable over metallic Li. However, not all Li coverage is removed, a dark deposit remains on the location of the original plated Li. The golden particles further change their color to brownish-red, indicating that they are not yet fully charged, and hence possess enough capacity to further intercalate Li. This supports the assumption that the dark deposit corresponds to irreversibly plated Li that will not be intercalated, even after extended relaxation periods. There are two possible explanations for the dark coloration; either the dark deposit consists of plated Li that lost electrical contact to the graphite during ongoing Li dissolution, or it consists of SEI-remnants of the plated Li on the surface; the exact origin cannot be determined. However as both effects lead to irreversible Li loss, both are undesirable in commercial cells.

A two-electrode cell was used for the optical detection in this work, instead of a three-electrode cell. Measurements have to be done using a reliable reference electrode, to measure the graphite potential without influence of the counter electrode. Also, this measurement is a representative measurement due to the size of the two/three electrode cell as opposed to the large EV cell. Therefore, it doesn’t account for the temperature distribution and delta temperature of the big cell. Another drawback of optical detection method is that they are generally limited to imaging surface only. It is difficult to arrive at a quantitative estimation.

Computed Tomography (CT)

K.J. Harry et al used synchrotron hard X-ray microtomography on symmetric lithium/polymer/lithium cells. The polymer electrolyte was a polystyrene-block-
poly(ethylene oxide) copolymer mixed with lithium bis(trifluoromethanesulfonyl)imide (LiTFSI) salt. Computed tomography (CT) measurements allow high resolution imaging and are not surface limited but need to be performed with an optimal radiation power to avoid radiation damage. Consequently, highly optimized synchrotron measurements are necessary for two- and three-dimensional imaging of delicately structured metallic Li. The cells were first cycled 15 times at a low current density of 0.02 mA cm\(^{-2}\), followed by cycling at a current density of 0.175 mA cm\(^{-2}\) until the experiment was stopped for cell imaging or when the cell shorted. The imaging was performed using monochromatic hard X-rays with energies chosen in the 22-25 keV range generated by the synchrotron. The X-ray shadow cast by the sample was converted into visible light using a scintillator. An optical microscope magnified this image and converted it into a digital image file. The sample was then rotated by a fraction of a degree and repeatedly imaged until 1,025 images were collected from the sample as it was rotated through 180°. The volume of dendrites was calculated by assuming each dendrite structure is composed of two half ellipsoid and measuring their length from the collected images. Their data show conclusively that under every dendritic structure there were subsurface structures located within the lithium electrode. In the early stages of dendrite formation, the volume occupied by the subsurface structure is significantly larger than that occupied by the dendritic structure protruding out from the electrode surface. This work indicates that preventing the growth of lithium dendrites may involve suppressing the nucleation of subsurface structures.

Later A. Yermukhambetova et al extended this method to explore the Li-S battery by 3D in situ X-ray tomography.\(^{[197]}\) They used a multi-scale, 3D X-ray imaging approach to examine an electrode both in situ at the micro-scale and ex situ at the nano-scale for a micron sized elemental sulfur and carbon black composite cathode. Sun et al used synchrotron in-line phase contrast X-ray CT to non-destructively visualize the evolution of electrodeposited lithium and observed separator rupture caused by their growth.\(^{[198]}\) O.O. Taiwo et al applied 3D X-ray CT to investigate the cycling graphite/Li half-cell in situ.\(^{[199]}\) Operando synchrotron X-ray CT enabled the real time 3D visualization of pit formation at the Li metal surface due to Li dissolution. Time-lapse laboratory X-ray CT imaging was used to track the growth of moss-like lithium deposits at the Li electrode surface over longer periods of cycling. The 3D imaging showed the appearance of pit-like holes on the Li metal surface as a result of Li dissolution during the first discharge, and the formation of a mossy, micro-structured lithium layer that increased in thickness with repeated cell cycling, penetrating the separator in the process. The results also highlight the importance of separator morphology in controlling or preventing dendrite growth in high tortuosity separators within commercial lithium batteries. One caution to be considered during measurement is the X-ray radiation dose to the cell samples must be carefully chosen to prevent material degradation and cell performance losses.

Recently R.F. Ziesche et al studied 4D imaging of primary lithium-batteries (Li\(_x\)MnO\(_2\)) using correlative neutron and X-ray tomography.\(^{[200]}\) The non-destructive nature of both X-ray and neutron CT provides opportunities for four-dimensional (4D) studies to explore the evolution of three-dimensional (3D) structures with time. After a series of data processing steps, these shadow images were converted to cross-sectional slices that were then stacked together to render a 3D reconstruction of the cell. Highly neutron-attenuating elements, such as Li and hydrogen, appear bright in the neutron tomograms, which facilitates detection of the Li anode between the darker cathodes of
lower attenuation. X-ray CT allowed for the quantification of mechanical degradation effects such as electrode cracking from the electrode bending process during cell manufacturing. Neutron CT gives information about the electrochemistry such as Li-ion transport and intercalation, electrolyte wetting of the cell and consumption or gas formation. The bulk properties that computational models were previously based on may now be replaced by spatially resolved, transient values. This technique offers the capability of incorporating local heterogeneities and would be able to provide a more accurate and comprehensive description of the practical operation, degradation and failure of a cell. Due to the electrolyte and Li having similar attenuation coefficients for neutrons, a direct correlation between the SOC and the amount of Li consumed could not be made, highlighting a limitation of this technique for the most commercially relevant cells.

Nuclear Magnetic Resonance Spectroscopy

Li NMR spectroscopy is a noninvasive method for investigating the structural changes that can occur in electrode materials, when performed in situ.\[^{[501-204]}\] The \(^{7}\text{Li}\) signal (spin = 3/2, 92.5% abundance) can be acquired on a time scale that is much faster than the typical charge-discharge cycle. \(^{7}\text{Li}\) NMR shows a chemical shift at 245±1 ppm for pristine Li metal and at >258 ppm for plated Li. The structural changes that occur in the active material at various states-of-charge can be detected and quantified by recording spectra at suitable time intervals. Nuclear magnetic resonance imaging (MRI) has proven to be capable of Li microstructure imaging, either directly or indirectly.

Hsei et al studied \(^{7}\text{NMR}\) on Li symmetric cell by both ex situ and in situ NMR\[^{[205]}\] to quantify irreversible lithium losses in batteries. Using NMR, GC-MS and calculations they distinguished losses due to SEI formation and fractions of dead lithium. Revealing a distribution of different lithium metal microstructures on both working and counter electrodes upon plating and stripping. Using these combination techniques, they estimated dead lithium fractions of 3.3% ± 0.6% (with 5% FEC) and 9.4% ± 0.6% (without 5% FEC). NMR as a standalone technique can suffer from issues of sensitivity and selectivity. Its resolution is lower than with optical or electron microscopy techniques. In fact, resolutions below 100 µm are hard to obtain for Li.\[^{[186]}\] NMR method does not directly distinguish between the mossy and the dendritic Li structures.

Dynamic nuclear polarization (DNP) is another promising approach that uses ~10\(^3\) times greater gyromagnetic ratio of paramagnetic electrons to hyperpolarize nuclear spins. In typical exogenous DNP experiments, organic radicals are added to the sample and it requires cryogenic sample cooling. The SEI on reduced graphene oxide and silicon anodes has been studied by exogenous DNP, whereby a solution of organic radicals is added to the system before cooling to 100 K or below, slowing the electron relaxation times of the radicals so that the ESR transition can be more easily saturated.\[^{[206-208]}\] This method has several limitations related to voltage instability of organic radicals, alteration of SEI by radicals, and it is not selective to SEI. Recently Hope et al achieved hyper polarization by irradiation of the conduction ESR (CESR) transition and partially equalizing the spin band. The electrons then cross relax with \(^{7}\text{Li}\) nuclei in the metal, thereby inducing nuclear hyperpolarization.\[^{[209]}\] Selective enhancement of the SEI was observed in the diamagnetic \(^{7}\text{Li}\), \(^{1}\text{H}\) and \(^{19}\text{F}\) NMR spectra, and the relative DNP enhancements allow the proximity of
different species to the metal surface to be inferred. By using double resonance experiments they identified the components of the SEI such as polymeric organic species and LiF. Effect of electrolyte was evaluated by using electrolytes with and without the common additive fluoroethylene carbonate (FEC). Their experiments showed that appreciable \(^7\text{Li}\) enhancements can still be achieved by Overhauser DNP under static conditions.

**Neutron Diffraction**

Neutron diffraction offers a promising alternative to derive information about the processes within a Li-ion battery. In contrast to x-ray diffraction, it has the advantage of being sensitive to lighter elements (like Li, oxygen and nitrogen) and the high penetration depth of thermal neutrons makes it possible to investigate Li-ion cells in a non-destructive way. There are two approaches to neutron diffraction on Li-ion batteries: one is to build special *\textit{in situ}* cells, optimized for the diffraction experiment (special geometry, deuterated electrolyte, etc.) \cite{25, 26, 210, 211}. The other approach is to use commercial cells, with the benefit that electrochemical performance and balancing of the electrodes are optimized by the manufacturer. The detail method and instrument specification can be found in Ref. \cite{40}.

Metallic Li in contact with the graphite anode is unstable; Li reacts with the graphite and diffuses into its interstitial spaces during intercalation. Even at -20°C, most of the metallic Li is intercalated after only 13 h. This makes conventional studies on Li plating difficult: since the Li diffusion coefficient at 23°C is reported to be approximately six times the diffusion coefficient at -15°C \cite{38, 212}, at room temperature at most possibly less Li remain to disassemble a cell or study Li plating by other methods. This fact further emphasizes the importance of using *\textit{in situ}* techniques when Li plating is studied. Valuable insights into Li plating in commercial cells can be obtained by neutron diffraction.

![Figure 6: Diffraction data collected during C/5 charge at -20°C (CV phase begins after 57.3% SOC); data collected every 30 min are shown, SOC relative to nominal RT](https://mc04.manuscriptcentral.com/jes-ecs)
Zinth et al investigated the transformation of LiC\textsubscript{12} to LiC\textsubscript{6} during the charging process in detail\cite{Zinth}. Intensities of the signal reflection of each of the two phases were analyzed, as described in Figure 6. It was shown that the loss of cyclable Li was significantly increased for higher charge currents, when cells were operated at sub-ambient temperatures of -20 °C, indicating that a considerable amount of Li has been deposited on the anode's surface. During a following rest period, the signal intensities of both phases changed, indicating a continuous intercalation of Li into the host structure, even after the applied current has been switched off\cite{Zinth}.

When the graphite anode is lithiated to more than 50%, the two existing phases LiC\textsubscript{6} and LiC\textsubscript{12}, commonly denoted as stage 1 and stage 2 in literature, show characteristic reflections at scattering angles 2\theta for the LiC\textsubscript{6} (001) at 33.53° and for the LiC\textsubscript{12} (002) at 35.04°. The intensity of the Bragg reflections increases with increasing share of the respective phase. Therefore, diffraction data in the range of 30°-40° 2\theta were collected in 2 min intervals throughout the experiment to follow simultaneously the strong (002) reflection of graphite and the similarly prominent reflections of the intercalation compounds LiC\textsubscript{12} and LiC\textsubscript{6} during the charging and relaxation. The integral intensity of the reflections was extracted by pseudo Voigt fitting of the data. In case of overlapping reflections, the full width at half maximum (FWHM) was refined once for all reflections.

The plated Li was compared with the voltage relaxation peaks. The changing peak intensities match distinctly with the voltage relaxation behavior, measured at the cell terminal. It can be observed that the amount of plating Li correlates non-linearly with the charging rate, whereas a charging current of C/2 leads to a deposited amount of Li of 5.5% and a current of 1C to 9.0% of the charged capacity. Above a threshold current of C/5, Li plating increases dramatically. The results indicate that the intercalation rate of reversibly deposited Li seems to be constant, independent of the deposited amount, and that the duration of the voltage plateau during the relaxation period can be used as an indicator for deposited Li.

**Electron Paramagnetic Resonance**

Electron Paramagnetic Resonance Spectroscopy (EPR), is a non-invasive and highly specific technique capable of detecting unpaired electrons or radicals. Though EPR has been successfully used in other fields it gained interest in the LiB field recently. It is a powerful tool for characterizing the formation and disappearance of radical oxygen species during the cycling of a battery. EPR offers unprecedented opportunities to study high-capacity electrode materials. With EPR imaging (EPRi), it is possible to view the electron-density distribution in materials, if we can develop EPR imaging (EPRi). The electrochemical cell for the study should be made up of a material that is (i) compatible with all cell components (electrolytes, electrodes and so on), (ii) transparent to microwave radiation, and (iii) inactive for EPR.
Electron paramagnetic resonance (EPR) spectroscopy is more sensitive than nuclear magnetic resonance, while still employing low energy radiation that does not affect the chemical properties or the morphology of the investigated samples. Conduction EPR (CEPR) was carried more than 50 years ago for the detection of conduction electrons in metallic Li\textsuperscript{[213, 214]}. Recently there has been a growing interest for the in operando study of Li plating in a LiB\textsuperscript{[30]} and the quantitative analysis of Li plating on graphite\textsuperscript{[31]}, both are very challenging with alternative techniques. They constructed a special cell with graphite working electrode and a lithium iron phosphate (LFP)-based counter and reference electrode. The occurrence of a metallic lithium EPR signal during charge was directly assigned to lithium plating on the graphite electrode. A small quantity of manganese (II) oxide (MnO), was placed outside of the cell to serve as an inert EPR reference. Operando EPR allowed to deconvolute the coulombic inefficiency of the lithium plating/stripping process thereby identifying the contributions derived from dead lithium vs. SEI formation, which cannot be achieved with standard electro-chemical techniques. Furthermore, the EPR spectra also contain spatial information regarding the location of lithium metal deposits (separator vs. electrode).

Sathiya et al first designed a special EPR electrochemical cell for in situ characterization\textsuperscript{[215]} The cycling profile of Li\textsubscript{2}Ru\textsubscript{0.75}Sn\textsubscript{0.25}O\textsubscript{3} electrode versus Li was studied by collecting and analyzing the EPR image of the cathode material. In the cathode material Ru\textsuperscript{5+}, O\textsuperscript{2-} and O\textsuperscript{2-} are EPR active and Ru\textsuperscript{4+} and O\textsuperscript{22-} are EPR inactive. As the change in the oxidation state of cathode components changes their EPR activity, it is possible to monitor the charge/discharge cycling by EPR and EPR image. Similarly, by CEPR method, Li morphology could also be distinguished from the linewidth of the metallic Li signal, with porous Li showing a much narrower EPR line than bulk Li, as can be seen in Figure 7. Therefore, CEPR is a sensitive indicator for the dimensions of the Li structures.

Neimöller et al used CEPR to study dendrite growth on separator\textsuperscript{[216]} They explained how the EPR line shape serves to distinguish different kinds of Li morphologies. EPR of conductive samples are influenced by the skin effect, which leads to a depth dependent attenuation and sample thickness dependent phase shift of the EPR signal\textsuperscript{[215, 217]} The skin depth is about \(\delta \approx 1.1 \, \mu m\) for metallic Li at X-band microwave frequencies. When the sample thickness \(d\) is larger than the skin depth \(\delta\), the experimentally acquired first derivative EPR spectrum shows a Dysonian lineshape that can be approximated by a phase-shifted Lorentzian\textsuperscript{[213]} The resolution achievable in an EPR experiment is similar to that of MRI\textsuperscript{[218, 219]}, which depend on the full width at half maximum of the point spread function\textsuperscript{[220]}. The EPR linewidth in the absence of a magnetic field gradient and the applied imaging gradient determine the image resolution. The pixel or voxel length \(\Delta z\) of the EPR image is proportional to the EPR peak-to-peak linewidth \(\Delta H_{pp}\) and to the inverse of the gradient \(G\textsuperscript{[221]}\), narrower lines and stronger gradients allow for higher resolution images.

\[
\Delta z \propto \frac{\Delta H_{pp}}{G} \tag{55}
\]
A. Figure 7: First derivative CEPR signal, as measured in a field swept EPR experiment, for metallic Li with different morphologies. The peak-to-peak linewidth for dendritic Li with ca. 0.005 mT (green), mossy Li is 0.03 mT (blue) and bulk Li is ca. 0.15 mT (red). Reproduced from Neimöller et al Scientific Reports, 8, 14331 (2018).[216]. (CC BY, https://creativecommons.org/licenses/by/4.0)

The self-diffusion of spin-carrying conduction electrons changes the linewidth. As the spin depth is much bigger than the skin depth, electrons diffuse from the surface into the bulk of the material. As a result, excited electron spins are withdrawn from a region excited by the external microwave field and replaced by electrons with an equilibrium spin state[214]. This process leads to a line broadening. Therefore, when a narrow dendritic line is observed, all the spins are contained in structures that are fully penetrated by the microwave field, hence quantitative imaging of the spin density is possible. This imaging method shows great potential for imaging of Li dendrite growth within solid electrolytes. Quantitative in situ application of CEPRI faces a challenge of multiple peaks from different Li morphologies and other paramagnetic species. As mentioned earlier, the other challenge with EPR analysis is that it requires the construction of a special cell which is transparent but inactive to EPR.

Accelerated Rate Calorimetry (ARC)

Similar to chemical reactions, electrochemical reactions have thermal signatures[222, 223]. Sample temperature variations alter the kinetics of redox processes and material properties that influence mass transport such as electrolyte viscosity. These temperature variations can permit undesirable outcome such as Li plating and excessive solid-electrolyte interphase (SEI) growth. Low temperatures, high charging rates (C-rates), and spatial thermal gradients are known to enable Li plating[46]. High temperatures accelerate SEI growth and diminish performance. On the cell level, a measure for cell safety is the onset temperature of exothermic reactions in accelerated rate calorimetry (ARC) tests. In
ARC tests, cells are heated step-by-step under quasi-adiabatic conditions until exothermic reactions are detected (heat-wait-seek experiment).

Thomas Waldmann and Margret Wohlfahrt-Mehrens studied the Li plating in a commercial NCA/graphite based 18650 cells after cycling at 0°C. They attributed the thermal runaway in their study to the stronger exothermic reactions of metallic Li. A more violent decomposition was noticed in the Li plated cells in the ARC tests. The damage pattern reproducibly included ejection of the jelly roll from the cell housing. This damage pattern is not observed for fresh cells without Li plating. The fresh cells without Li plating took the longest time to reach the maximum temperature of thermal runaway (1950 min). All tested cells with Li plating showed significantly shorter times to reach the thermal runaway and the maximum temperatures. Their study demonstrated the influence of time interval between cycling and ARC measurement on the onset of thermal runaway. For rest times of 1.5 h and 8d after Li plating, the maximum temperatures were reached after 225 min and 1000 min, respectively. Carter et al. evaluated three fully charged cells by ARC to assess their thermal runaway. One cell was conventionally cycled equilibrium 20°C and it began to self-heat at 130°C. The second cell was cycled at equilibrium 0°C and the third was cycled with transient 40 to 0°C. The latter two cells began to self-heat at low temperatures 50 and 35°C, respectively. The higher peak temperature of the transient 40 to 0°C cell (154°C compared to 108°C for the Equilibrium 0°C condition) caused by Li side reactions, indicates a greater quantity of metallic Li, augmenting the heat generation. Since venting is not evident in the data for each of these cells, both cans rupture during the thermal runaway event, releasing active material. In the case of the equilibrium 0°C cell, the jellyroll ejects through the top of the cell whereas the transient 40 to 0°C cell exhibits side wall rupture opening the cell in the radial direction. In contrary, the conventional 20°C cycled cell didn’t have any rupture, although it also combusted after 200°C.

Charging Li-ion cells under spatially uniform, temporally transient thermal conditions can greatly reduce performance and safety. A single charging of a commercial Li-ion battery while its temperature is decreasing temporally toward 0°C, is enough to cause Li plating and early loss of capacity, earlier onset temperature of thermal runaway, and higher heat generation temperature during thermal runaway. Since cells charged during the transient to 0°C experience significant Li plating and subsequent stripping, its long-term performance differed from a cell initially equilibrated to the same condition. The equilibrium 0°C cell undergoes significantly less Li metal plating and stripping (~0.2 Ah) but eventually experienced jellyroll collapse after ~30 cycles due to low temperature. However, the largest quantity of Li plated and stripped in the transient 40 to 0°C case (1.04 Ah) which caused the jellyroll to collapse more rapidly, after only ~5 cycles. These distinct degradation modes enable extremely energetic failure modes following self-heating at mild temperatures of only 50 and 35°C for the equilibrium 0°C and transient 40 to 0°C cases, respectively. The Transient 40 to 0°C case shows evidence of more residual Li metal and a more rapid acceleration to thermal runaway leading to can wall breaches.
The cells’ sensitivity to thermal transient conditions emphasize the need to understand the correlations between Li-ion battery environments and local electrochemistry. This understanding can better evaluate safe operating conditions and thermal management strategies. The thermal sensitivity of form factors and pack assemblies should be evaluated. Moreover, it is important to understand the time to equilibrium when selecting C-rate, so that substantial SOC variation does not occur during a thermal transient condition. These considerations are particularly important for extreme fast charging. One main drawback of the ARC study is that the instrument is commonly designed to study small cells, such as 18650 or 26650 sizes. In order to study bigger EV cells, it must be custom designed for safe measurement which increases the cost.

In Operando Acoustic Detection of Li Metal Plating

Non-destructive methods for detecting localized Li metal plating in LiBs have been challenging to set up. An elegant and straightforward method is to correlate cell thickness measurements with plating in operando[225, 226]. Although this is a proven method, it is unable to distinguish between gassing reactions and Li metal plating, which may both increase cell thickness. Earlier report by Hsieh et al, introduced the ability of acoustic ultrasound to probe physical dynamics of a battery during cycling[227]. Specifically, Davies et al[228] attributed the shifts in the transmitted ultrasonic signals due to variations of the electrode density and bulk modulus. Similarly, Gold et al.[229] utilized a comparably lower frequency ultrasonic pulse (200 KHz versus 2.25 MHz) to analyze the arrival time of the slow, compressional wave. J.B. Robinson et al studied a commercial cell phone battery by successfully employing spatially resolved ultrasound acoustic measurement combined with X-ray computed tomography.[230] They demonstrated the applicability of this method for battery state of health measurement and also identified local cell construction glitches associated with anode tab. Z. Deng et al along with Jeff Dahn have successfully demonstrated cell wettability differences in pouch cells by employing ultrasonic scanning technology[231].

Bommier et al.[232] demonstrated that the simple and scalable in operando acoustic technique can be utilized to detect Li metal plating in commercial scale LiBs. The detection method is semiquantitative, as more significant Li metal plating results in increased hysteresis of the acoustic time of flight. The acoustic technique can observe localized and heterogeneous plating events, because the signal comes from the total acoustic pathway between the locations on the cell connected to the transducers. Moreover, the technique can be deployed with minimal start-up costs, requiring transducers, an acoustic detector, and an open source software package.

Bommier et al studied the cycling of the LiCoO2/graphite battery cell. The first two C/15 cycles show a clear Time Of Flight (TOF) shift, with the TOF decreasing during the charging step and increasing during the discharge step. Graphite has a larger volumetric-specific capacity than LiCoO2 (LCO), therefore it takes up a larger portion of the acoustic pathway and will have a greater impact on the TOF. The local inflection points in the TOF shift may be due to graphite staging events, as these are located at the same SOC as inflection points in the voltage curve. The TOF shifts for the initial slow C/15 charge and the faster fixed capacity charge were compiled for all tests and analyzed. As described in Figure 8, the acoustic TOF shift between the C/15 current rate (blue curve) and the 0.5C
fixed capacity charge (red curve) is similar. However, as the current rate is increased to 0.75C and 0.88C, the respective TOF shift between the slow charge and the faster fixed capacity charge begins to deviate (Figures 8B and 8C). At the current rate of 1C, there is a clear anomaly that develops within the cell at around 0.070 Ah: the initially decreasing TOF shift suddenly reverses course and starts to increase. When the fixed capacity charge is finished, the TOF shift is completely different from that of the slow rate C/15 charge, with a large difference in the endpoints (Figure 8D). These results indicate the possibility of using the TOF endpoint difference as a plating detection parameter.

Figure 8. Comparison of acoustic TOF shifts during the C/15 charge and the fixed-capacity charge of 0.210 Ah, at four different current rates. No voltage cutoffs were set for fixed-capacity charges. These cells were cycled at 10°C with a fixed capacity charge of (A) 0.5C, (B) 0.75C, (C) 0.875C, and (D) 1C. Reproduced from C. Bommier et al, Cell Reports Physical Science 1, 100035, (2020)[232]. (CC BY, https://creativecommons.org/licenses/by/4.0)
Figure 9: (A) Plot of logarithmic rate (1/t) versus inverse temperature (1,000/T),
using the associated color bar to factor in the time of flight (TOF) endpoint
difference between the C/15 charge and the fixed capacity charge. (B) Optical
images of a delithiated graphite anode control (left of image) and the anode at the
corresponding fixed-capacity charge (right of image). The colored dot represents
the TOF endpoint difference according to the color bar used in (A)\[232]\.
Reproduced from C. Bommier et al, Cell Reports Physical Science 1, 100035,
(2020). (CC BY, https://creativecommons.org/licenses/by/4.0)

Figure 9 shows the plot of logarithmic rate versus inverse temperature. A symmetric
logarithmic color bar has been used for easier identification of TOF endpoint differences.
Lower TOF endpoint differences are indicated by an increasingly dark green color, and
higher TOF endpoint differences are indicated by an increasingly dark magenta color. The
depicted equation represents the linear fit (Arrhenius relationship) between logarithmic rate
and inverse temperature, shown by the dotted line. The gray error bar represents the
statistical standard error, which is the variance of the dataset divided by the set size. TOF
endpoint differences generally increase with rate at any given temperature, as visualized
by the color of the datapoint. Because plating is a continuum event, the linear fit generally
describes conditions where plating is likely (i.e., one standard deviation above the linear
fit) but does not conclude with certainty the presence of plating due to the statistical nature
of this experiment. With a pouch cell that had undergone visible gassing and swelling, the
permanent attenuation of the acoustic signal can be observed. Acoustic waveforms are
known to propagate much more easily through solid or liquid mediums due to a lower
impedance mismatch ratio, and significant attenuation of the signal implies the presence of
gassing.

One drawback of the technique as it currently stands is the bulk methodology. A simple
ultrasonic pulse through the length of the battery allows for a quick assessment of the entire
stack. At the same time, it also convolutes contributions from individual electrodes.
Therefore, it would be difficult to distinguish whether plating occurred on only one or
several electrodes and determine which electrodes plated. The acoustic effects are solely
based on the area between the transducers, so to identify spatial heterogeneities, multiple
transducers can be placed on a single cell, given that it is large enough. This makes for
useful commercial cell quality analysis, especially owing to the low cost and speed at which
diagnostics can be carried out. As the wavelength of the ultrasonic pulse is relatively low,
on the order of ~1 mm, there may exist a minimum resolution beyond which plating cannot
be detected (e.g., single-layer electrode pouch cells). However, W. Chang et al
demonstrated the sensitivity of thickness detection can be improved to micron resolution
when combined with in-situ transmission X-ray microscopy (TXM)\[233]\. They combined
the ultrasound acoustic detection with TXM to study the expansion and contraction of a
pouch cell during cycling. Their expectation is that TXM has sufficient range and pixel
resolution to measure both the total cell thickness and the average layer thicknesses, though
it proved to be more difficult to image. Nevertheless, they successfully measured the
effective stiffness of a commercial pouch cell battery during cycling from the wave velocity
and cell thickness in operando by the TXM/acoustics setup. Note that cell stiffness is
resulting from electrode expansion during cycling. Lastly, the acoustic detection sensitivity
is also related to the attenuating properties of the media. Form factors that include materials
that have higher impedance mismatch with the transducers or that intensify waveform convolution (e.g., springs/spacers in coin cells) may be more difficult to analyze.

Mass Spectrometry Titration

Recently Mass spectrometry titration (MST) was successfully employed for inactive Li quantification on graphite with a detection limit of 20 nmol Li, which corresponds to 0.5 μAh of total plated Li. This technique is a modified version of a similar titration to quantify Li₂CO₃ on extracted Li-O₂ battery cathodes and was recently used to quantify surface Li₂CO₃ on Li-rich and Li-stoichiometric NMC cathodes. With the MST, it was possible to precisely identify: 1) the charge rate at which Li plating is observed if the electrode is charged to its nominal full capacity (corresponding to 372 mAh/g graphite), and 2) the state of charge at which Li plating commences at a 4C rate (full charge in 15 minutes). It can be characterized by an electrochemical model to determine the Li plating exchange current density and the efficiency of plated Li stripping during discharge.

M.J. Shane et al, used MST to quantify plated and inactive Li. The graphite electrode sample was first placed in an airtight titration vessel, the apparatus and the method are described previously. Previous work on quantification of Li₂CO₃ on cathode materials using the same MST setup showed that any concentration above 1M H₂SO₄ provided sufficient excess acid to react all present carbonate species, and H₂SO₄ was chosen over other acids to remain consistent with similar previous carbonate titration studies. Carbonate titrations using phosphoric acid have also been reported and yield similar results. The maximum amount of gas evolved in any single experiment was roughly 15 μmol, which is calculated to cause a pressure increase in their sealed titration vessel ~30 torr, not large enough to break the titration vessel. An inert gas should be used as the carrier gas to ensure evolved H₂ does not combust given the high exothermicity of Li reacting with acid or water.
As seen in Figure 10, it is possible to achieve a detection limit of 20 nmol inactive Li with MST, which is based upon the expected minimum H$_2$ signal that can be quantified upon titration. One mole H$_2$ corresponds to 2 moles of inactive Li. The lowest amount of inactive Li measured with MST was 105 nmol Li, which is much higher than the signal to noise.

The capacity loss due to SEI components such as carbonates and Li$_2$C$_2$ was calculated from the amount of CO$_2$ and C$_2$H$_2$ evolved whereas inactive Li was calculated from amount of evolved H$_2$. Inactive Li includes only electronically isolated Li and Li$_x$C$_6$. It can be calculated by assuming that plated Li is formed via a 1 electron/Li process and Li$_x$C$_6$ is formed via a x electron/Li$_x$C$_6$ process, i.e., 1 mole electron per mole inactive Li. This calculation is shown below, where $Q_{Li}$ is the inactive Li equivalent capacity and $n_{Li}$ is the number of moles of inactive Li measured via H$_2$ evolution upon titration:

$$Q_{Li} = n_{Li} \frac{1 \text{ mol e}^-}{1 \text{ mol Li}} \frac{96485 \text{ C mAh}}{1 \text{ mol e}^-} \frac{3.6 \text{ C}}{1 \text{ mol Li}}$$

[56]

Contribution of each source of capacity loss is shown in Figure 11. Solid carbonate and Li$_2$C$_2$ contribution to total irreversible capacity are calculated based on the increase in their amount above the baseline values observed after formation cycling. In order to estimate the contribution of Li dislodged from graphite and entrained in the separator, they constructed a cell with two Whatman QMA separators and titrated both the graphite electrode and the separator adjacent to it without rinsing. The separator adjacent to the Li metal counter electrode was excluded from the titration as it would have entrapped the Li ion from the counter electrode. Thus, mass spectrometric titration is a powerful technique in distinguishing between inactive Li and other surface components of plated Li. However, this technique is ex situ, which require tearing down the cell for analysis. The accuracy of the estimation depends on the time between tear down and analysis as with other ex situ techniques.
Figure 11: (a) Cell cycling procedure: two C/10 formation cycles (0.010-1.5 V) followed by three cycles of varying charge rate (shown in red) and C/4 discharge to 1.5V with 30-minute OCV period between each charge/discharge. (b) Total contributions of inactive Li, solid carbonate, Li$_2$C$_2$ and other species to the observed irreversible capacity after three cycles of varying charge C-rates from (a)$^{[48]}$. Adapted with permission from E.J. McShane et al, ACS Energy Lett., \textbf{5} (6) 2045–2051 (2020).$^{[208]}$ Copyright © 2020 from American Chemical Society

Summary and Outlook

Anode Li plating is one of the main factors that result in rapid aging and safety issues of LiBs; addressing this issue is critical to design safer and more durable LiBs. This paper presents an overview of existing methods and recent advancements in the detection of anode Li plating. The review covers the underlying mechanisms of Li plating, diagnostic techniques, their benefits and the current status of development.

Li deposition is more likely to happen with decreasing temperature, increasing SOC and increasing charging C-rate. However, the interplay of these main factors is also important. Low temperatures do not lead to Li deposition in all cases, however, if certain limits of SOC and/or charging C-rate are exceeded, Li deposition is seen to occur. For example, a cell might be charged by high currents at low SOC and later by lower currents without Li deposition. At low temperatures, the Li deposition might be suppressed by low charging currents or by preheating of batteries. Variations can be observed when inhomogeneities/gradients are present inside the cells: temperatures, SOCs and current densities. Such variations can lead to a variety of macroscopic morphologies for Li deposition, including homogeneous Li deposition (Li plating), local Li deposition and marginal Li deposition. As an example, local formation of gas bubbles or SEI growth and clogging of pores at high temperatures can lead to local Li deposition during cycling at 45 °C.

Besides these evidences for Li metal deposition in Li-ion cells, a variety non-destructive and destructive methods exist which are based on ageing behavior, analysis of electrochemical data, cells, and electrodes. In addition to effects on ageing, Li deposition is also connected to safety concerns. Two main safety mechanisms were identified including, internal shorts caused by the dendritic growth of Li and exothermic reactions of Li with electrolyte. The dendritic growth of Li is connected to the Li/Li$^+$ interaction on
the atomic level, the presence of ridges, and higher field densities at the tips of the dendrites. Besides the very low probability of failures due to Li dendrites in Li-ion cells (ppm range) they are relevant, and some attempts have been made for failsafe designs and early detection of dendrites. Exothermic reactions of Li metal with electrolyte can already start at 30–35 °C under the adiabatic conditions of ARC tests. In an application this translates to a situation, when not enough heat can be transferred from the cell to the surroundings, i.e. for a very high thermal isolation. This shows the importance of cooling systems for batteries in applications.

It is challenging to sensitively and quantitatively detect anode Li plating in practical LiBs, which is significant for battery management systems. The cause of anode Li plating depends strongly on the cell design and operating conditions, especially charge rate and temperature. During high rate charging at room temperature, the poor solid phase diffusion of Li ions is responsible for anode Li plating. As the temperature decreases, the interface reaction is also hindered in addition to solid phase Li diffusion. In order to suppress anode Li plating, optimizing the electrolyte composition and modifying the graphite surface structure by methods such as coating, and doping are usually employed. In addition, appropriate working conditions and charging protocols can reduce the possibility of anode Li plating. A recent report from Yang et al[238] described the performance of self-heating LiB structure which has thin nickel (Ni) foils embedded inside a cell that can create immense and uniform heating. The Ni foil is an inherent component of a single cell along with electrodes and electrolyte. It serves as an internal heating element, as well as an internal temperature sensor as its electrical resistance varies linearly with temperature. The Ni foils only adds 0.5% weight and 0.04% cost to a conventional LiB single cell. They were able demonstrate fast charging without Li plating at temperatures as low as -50°C with the help of the self-heating element.

Various techniques have been applied to detect anode Li plating. Among these, physical characterizations such as Conduction EPR, neutron diffraction and mass spectrometric titration are promising for understanding the mechanisms of Li plating, while electrochemical methods such as coulombic efficiency, discharge voltage profiles, electrochemical impedance and cell thickness analysis are more applicable to practical LiBs.

The following table provides the Li detection methods, their advantages and current challenges:

| Methods and advantages | Challenges |
|------------------------|------------|
| Three-electrode diagnostics | Temperature distribution differences between small 3 electrode cell and multi stack EV battery is not linear. |
| Detection of voltage plateau (dV/dt) | Interferences from graphite staging with Li plating at low charge rates. |
intercalation leads to a voltage plateau and can be identified in the plot of dV/dt vs %SOC. The most common method that can be studied on full cell as opposed to constructing a special cell.

| Measurement of coulombic efficiency (CE) | High precision coulometer and temperature chamber are needed for accurate diagnosis. |
|----------------------------------------|----------------------------------------------------------------------------------|
| When Li plating happens, coulombic efficiency decreases. Easier to implement on the cell level. | SEI decomposition reactions in the cell can lead to gassing which can affect the measurement accuracy. It is also influenced by the thickness measurement jig and measurement point. |
| Measurement of cell thickness | It requires step charging protocol like HPPC with rest period included during which cell impedance measurement is performed. |
| 1 Ah of plated Li causes a volume change of 0.49 cm$^3$, compared to 0.12 cm$^3$ for intercalated Li. Thickness changes from Li deposition reflect on cell thickness. Easier to implement on cell level. | SEI decomposition reactions in the cell can lead to gassing which can affect the measurement accuracy. It is also influenced by the thickness measurement jig and measurement point. |
| Cell Impedance Analysis | It requires step charging protocol like HPPC with rest period included during which cell impedance measurement is performed. |
| Charge transfer region of cell EIS changes with SOC. Li plating changes the trend from a good cell vs plated cell. Study can be performed on full cell. | Early exothermic reactions in ARC tests needs massive Li plating. ARC chamber must be custom made to test the large EV batteries. |
| Accelerated Rate Calorimetry (ARC) | Optical Detection SEM/TEM/Raman/XPS Plated cell is teared down and anode was analyzed by the above techniques for Li plating. |
| Li plating is associated with thermal signature which can be measured by ARC. It can be performed on full cell. | Optical Detection SEM/TEM/Raman/XPS Plated cell is teared down and anode was analyzed by the above techniques for Li plating. |
| Li Nuclear magnetic resonance (NMR) | Optical Detection SEM/TEM/Raman/XPS Plated cell is teared down and anode was analyzed by the above techniques for Li plating. |
| The chemical shift of Li metal in NMR spectra (245–270 ppm, depending on Li metal structure). The peak intensity in NMR spectra is proportional to Li amount and is used to quantify Li plating. | As in the case with optical detection methods, NMR gives direct evidence only microscopic and electrode level Li plating. It is not a cell level analysis. |
| **Electron paramagnetic resonance (EPR)** | It is in the early stage of development. |
|----------------------------------------|---------------------------------------|
| The EPR signal of Li metal is much narrower and its signal center is slightly upfield compared with the signal from the intercalated Li (Li_xC_6). The magnitude of EPR signal is related to Li metal amount. It is more sensitive than NMR. | |

| **X-ray and Neutron Computed Tomography** | It has the potential for real time battery internal component imaging combined with surface characterization. It is in the early stage of implementation. |
|------------------------------------------|----------------------------------------------------------------------------------------------------------------------------------|
| X-ray CT allowed for the quantification of mechanical degradation whereas Neutron diffraction can identify Li-ion transport and intercalation, electrolyte wetting of the cell. | |

| **Acoustic detection** | It is a bulk technique that can be performed on a full cell. It can identify Li plating but cannot identify the location of Li plating initiation. It can be combined with other optical detection methods to get both cell level and microscopic information. |
|------------------------|--------------------------------------------------------------------------------------------------------------------------------------|
| The transmitted ultrasonic signals shift due to variations of the electrode density and bulk modulus. Can be studied on full cell. Easy to implement. | |

| **Mass spectrometric titration** | It is a quantitative technique. It is a component level analysis and is also an **ex situ** technique. Therefore, it has the same challenges as those of other ex-situ techniques. |
|---------------------------------|----------------------------------------------------------------------------------------------------------------------------------|
| Titration of anode and estimation of off gasses by mass spectrometry. It can identify the difference between inactive Li and the SEI components on plated Li down to 20nmol concentration. | |

### Acknowledgments

The authors appreciate the discussions and feedback from Dr. Dan Steingart of Columbia University and Dr. Bryan D. McCloskey of University of California, Berkeley.

### References

1. C.-Y. Wang, G. Zhang, S. Ge, T. Xu, Y. Ji, X.-G. Yang and Y. Leng, *Nature*, **529**, 515 (2016).
2. V. Etacheri, R. Marom, R. Elazari, G. Salitra and D. Aurbach, Energy & Environmental Science, 4, 3243 (2011).
3. H. Jiang, Z. Lu, M. Wu, F. Ciucci and T. Zhao, Nano Energy, 23, 97 (2016).
4. C. R. Birkl, M. R. Roberts, E. McTurk, P. G. Bruce and D. A. Howey, J Power Sources, 341, 373 (2017).
5. M. Dubarry, G. Baure and A. Devie, J Electrochem Soc, 165, A773 (2018).
6. R. Xiong, Y. Duan, J. Cao and Q. Yu, Appl Energ, 217, 153 (2018).
7. S. J. Harris, E. K. Rahani and V. B. Shenoy, J Electrochem Soc, 159, A1501 (2012).
8. S. J. Harris, A. Timmons, D. R. Baker and C. Monroe, Chemical Physics Letters, 485, 265 (2010).
31. J. Wandt, P. Jakes, J. Granwehr, R.-A. Eichel and H. A. Gasteiger, *Materials Today*, 21, 231 (2018).
32. J. Arai, Y. Okada, T. Sugiyama, M. Izuka, K. Gotoh and K. Takeda, *J Electrochem Soc*, 162, A952 (2015).
33. H. Ge, T. Aoki, N. Ikeda, S. Suga, T. Isobe, Z. Li, Y. Tabuchi and J. Zhang, *J Electrochem Soc*, 164, A1050 (2017).
34. J. Burns, D. Stevens and J. Dahn, *J Electrochem Soc*, 162, A959 (2015).
35. Q. Liu, R. Petibon, C. Du and J. Dahn, *J Electrochem Soc*, 164, A1173 (2017).
36. B. Bitzer and A. Gruhle, *J Power Sources*, 262, 297 (2014).
37. B. Rieger, S. F. Schuster, S. V. Erhard, P. J. Osswald, A. Rheinfeld, C. Willmann and A. Jossen, *Journal of Energy Storage*, 8, 1 (2016).
38. C. Birkenmaier, B. Bitzer, M. Harzheim, A. Hintennach and T. Schleid, *J Electrochem Soc*, 162, A2646 (2015).
39. F. Grimsmann, T. Gerbert, F. Brauchle, A. Gruhle, J. Parisi and M. Knipper, *J Power Sources*, 365, 12 (2017).
40. V. Zinth, C. von Lüders, M. Hofmann, J. Hattendorff, I. Buchberger, S. Erhard, J. Rebello-Kornmeier, A. Jossen and R. Gilles, *J Power Sources*, 271, 152 (2014).
41. R. V. Bugga and M. C. Smart, *Ecs Transactions*, 25, 241 (2010).
42. M. Bauer, B. Rieger, S. Schindler, P. Keil, M. Wachtler, M. A. Danzer and A. Jossen, *Journal of Energy Storage*, 10, 1 (2017).
43. C. von Lüders, V. Zinth, S. V. Erhard, P. J. Osswald, M. Hofmann, R. Gilles and A. Jossen, *J Power Sources*, 342, 17 (2017).
44. T. Waldmann and M. Wohlfahrt-Mehrens, *Electrochim Acta*, 230, 454 (2017).
45. U. R. Koleti, T. Q. Dinh and J. Marco, *J Power Sources*, 451, 227798 (2020).
46. R. Carter, E. J. Klein, T. A. Kingston and C. T. Love, *Frontiers in Energy Research*, 7, 144 (2019).
47. M. T. Pham, J. J. Darst, D. P. Finegan, J. B. Robinson, T. M. Heenan, M. D. Kok, F. Iacoviello, R. Owen, W. Q. Walker and O. V. Magdysyuk, *J Power Sources*, 228039 (2020).
48. E. J. McShane, A. M. Colclasure, D. E. Brown, Z. M. Konz, K. Smith and B. D. McCloskey, *ACS Energy Letters* (2020).
49. W. Fang, O. J. Kwon and C. Y. Wang, *International journal of energy research*, 34, 107 (2010).
50. Y. Ji, Y. Zhang and C.-Y. Wang, *J Electrochem Soc*, 160, A636 (2013).
51. X.-G. Yang, C. Bauer and C.-Y. Wang, *J Power Sources*, 327, 414 (2016).
52. X.-G. Yang, G. Zhang and C.-Y. Wang, *J Power Sources*, 328, 203 (2016).
53. X.-G. Yang, T. Liu and C.-Y. Wang, *J Power Sources*, 342, 598 (2017).
54. L. Liu, J. Park, X. Lin, A. M. Sastry and W. Lu, *J Power Sources*, 268, 482 (2014).
55. L. Liu, P. Guan and C. Liu, *J Electrochem Soc*, 164, A3163 (2017).
56. B. Wu and W. Lu, *J Power Sources*, 360, 360 (2017).
57. P. Arora, M. Doyle and R. E. White, *J Electrochem Soc*, 146, 3543 (1999).
58. R. D. Perkins, A. V. Randall, X. Zhang and G. L. Plett, *J Power Sources*, 209, 318 (2012).
59. N. Legrand, B. Knosp, P. Desprez, F. Lapicque and S. Raël, *J Power Sources*, 245, 208 (2014).
60. X.-G. Yang, Y. Leng, G. Zhang, S. Ge and C.-Y. Wang, *J Power Sources*, 360, 28 (2017).
61. T. Danner, M. Singh, S. Hein, J. Kaiser, H. Hahn and A. Latz, *J Power Sources*, 334, 191 (2016).
62. S. Hein and A. Latz, *Electrochim Acta*, 201, 354 (2016).
63. M. A. Cabañero, J. Altmann, L. Gold, N. Boaretto, J. Müller, S. Hein, J. Zausch, J. Kallo and A. Latz, *Energy*, 171, 1217 (2019).
64. J. Feinauer, S. Hein, S. Rave, S. Schmidt, D. Westhoff, J. Zausch, O. Iliev, A. Latz, M. Ohlberger and V. Schmidt, *Journal of Computational Science*, 31, 172 (2019).
65. T. Waldmann, M. Wilka, M. Kasper, M. Fleischhammer and M. Wohlfahrt-Mehrens, *J Power Sources*, 262, 129 (2014).
66. J. Remmlinger, S. Tippmann, M. Buchholz and K. Dietmayer, *J Power Sources*, 254, 268 (2014).
67. M. Petzl, M. Kasper and M. A. Danzer, *J Power Sources*, 275, 799 (2015).
68. S. S. Zhang, K. Xu and T. Jow, *J Power Sources*, 160, 1349 (2006).
69. M. Ouyang, Z. Chu, L. Lu, J. Li, X. Han, X. Feng and G. Liu, *J Power Sources*, 286, 309 (2015).
70. R. C. Bhardwaj, T. Hwang and R. M. Mank, Modulated, temperature-based multi-CC-CV charging technique for Li-ion/Li-polymer batteries, in, Google Patents (2014).
71. Y. Itou and Y. Ukyo, *J Power Sources*, 146, 39 (2005).
72. R. C. Bhardwaj, T. Hwang and R. M. Mank, Modulated, temperature-based multi-CC-CV charging technique for Li-ion/Li-polymer batteries, in, Google Patents (2014).
73. T. R. Garrick, J. Gao, X. Yang and B. Koch, *Submitted to J. Electrochem. Soc.* (2020).
74. T. R. Garrick, J. Gao, X. Yang and B. Koch, *Submitted to J. Electrochem. Soc.* (2020).
75. J.-i. Yamaki, S.-i. Tobishima, K. Hayashi, K. Saito, Y. Nemoto and M. Arakawa, *J Power Sources*, 74, 219 (1998).
76. P. J. Mohr, B. N. Taylor and D. B. Newell, *Journal of Physical and Chemical Reference Data*, 80, 633 (2008).
77. S. Ahmed, I. Bloom, A. N. Jansen, T. Tanim, E. J. Dufek, A. Pesaran, A. Burnham, R. B. Carlson, F. Dias and K. Hardy, *J Power Sources*, 367, 250 (2017).
78. A. Smith, J. Burns and J. Dahn, *Electrochem Solid St*, 13, A177 (2010).
79. J. Burns, N. Sinha, G. Jain, H. Ye, C. M. VanElzen, W. Lamanna, A. Xiao, E. Scott, J. Choi and J. Dahn, *J Electrochem Soc*, 159, A1105 (2012).
80. K. Zachary M., M. Eric J. and M. Bryan D., *Voltage Relaxation to Detect the Onset of Lithium Plating on Graphite for Fast Charging* (2020).
81. X.-G. Yang, S. Ge, T. Liu, Y. Leng and C.-Y. Wang, *J Power Sources*, 395, 251 (2018).
82. C. Fear, T. Adhikary, R. Carter, A. N. Mistry, C. T. Love and P. P. Mukherjee, *ACS Applied Materials & Interfaces*, 12, 30438 (2020).
83. F. Ringbeck, C. Rahe, G. Fuchs and D. U. Sauer, *J Electrochem Soc*, 167, 090536 (2020).
84. S. E. O’Kane, I. D. Campbell, M. W. Marzook, G. J. Offer and M. Marinescu, *J Electrochem Soc*, 167, 090540 (2020).
85. I. D. Campbell, M. Marzook, M. Marinescu and G. J. Offer, *J Electrochem Soc*, 166, A725 (2019).
86. K. A. Smith, *IEEE Control Systems Magazine*, 30, 18 (2010).
87. A. Seaman, T.-S. Dao and J. McPhee, *J Power Sources*, 256, 410 (2014).
88. M. Doyle, T. F. Fuller and J. Newman, *J Electrochem Soc*, 140, 1526 (1993).
89. K. E. Thomas, J. Newman and R. M. Darling, in Advances in lithium-ion batteries, p. 345, Springer (2002).
90. K. E. Thomas and J. Newman, J Electrochem Soc, 150, A176 (2003).
91. J. S. Newman and C. W. Tobias, J Electrochem Soc, 109, 1183 (1962).
92. J. Newman and W. Tiedemann, AIChE Journal, 21, 25 (1975).
93. J. Newman and K. E. Thomas-Alyea, Electrochemical systems, John Wiley & Sons (2012).
94. J. S. Dunning, D. N. Bennion and J. Newman, J Electrochem Soc, 120, 906 (1973).
95. J. S. Dunning, D. N. Bennion and J. Newman, J Electrochem Soc, 118, 1251 (1971).
96. M. Tang, P. Albertus and J. Newman, J Electrochem Soc, 156, A390 (2009).
97. D. Ren, K. Smith, D. Guo, X. Han, X. Feng, L. Lu, M. Ouyang and J. Li, J Electrochem Soc, 165, A2167 (2018).
98. X. M. Liu, A. Fang, M. P. Haataja and C. B. Arnold, J Electrochem Soc, 165, A1147 (2018).
99. S. Szabó, International Reviews in Physical Chemistry, 10, 207 (1991).
100. K. Yan, Z. Lu, H.-W. Lee, F. Xiong, P.-C. Hsu, Y. Li, J. Zhao, S. Chu and Y. Cui, Nature Energy, 1, 1 (2016).
101. I. Morcos, Journal of Electroanalytical Chemistry and Interfacial Electrochemistry, 66, 250 (1975).
102. T. Tran and K. Kinoshita, Journal of Electroanalytical Chemistry, 386, 221 (1995).
103. L. Gasparotto, N. Borisenko, N. Bocchi, S. Z. El Abedin and F. Endres, Physical Chemistry Chemical Physics, 11, 11140 (2009).
104. C. A. Paddon and R. G. Compton, The Journal of Physical Chemistry C, 111, 9016 (2007).
105. Y. Mo, Y. Gofer, E. Hwang, Z. Wang and D. A. Scherson, Journal of Electroanalytical Chemistry, 409, 87 (1996).
106. M. L. Urquiza, M. Otero, G. L. Luque, D. Barraco and E. P. Leiva, Electrochim Acta, 208, 92 (2016).
107. M. K. Datta, B. Gattu, R. Kuruba, P. M. Shanthi and P. N. Kumta, J Power Sources, 467, 228243 (2020).
108. M. Verbrugge, D. Baker and X. Xiao, J Electrochem Soc, 163, A262 (2016).
109. M. Verbrugge, D. Baker, B. Koch, X. Xiao and W. Gu, J Electrochem Soc, 164, E3243 (2017).
110. D. R. Baker and M. W. Verbrugge, J Electrochem Soc, 165, A3952 (2018).
111. D. R. Baker, M. W. Verbrugge and W. Gu, J Electrochem Soc, 166, A521 (2019).
112. D. R. Baker and M. W. Verbrugge, J Electrochem Soc, 167, 013504 (2019).
113. H. Zhang and M.-Y. Chow, in IEEE PES General Meeting, p. 1 (2010).
114. S. Abu-Sharkh and D. Doerffel, J Power Sources, 130, 266 (2004).
115. R. C. Kröeze and P. T. Krein, in 2008 IEEE Power Electronics Specialists Conference, p. 1336 (2008).
116. M. Chen and G. A. Rincon-Mora, IEEE transactions on energy conversion, 21, 504 (2006).
117. L. Gao, S. Liu and R. A. Dougal, IEEE transactions on components and packaging technologies, 25, 495 (2002).
118. F. Baronti, G. Fantechi, E. Leonardi, R. Roncella and R. Saletti, in IECON 2010-36th Annual Conference on IEEE Industrial Electronics Society, p. 2329 (2010).
119. O. Erdinc, B. Vural and M. Uzunoglu, in 2009 International Conference on Clean Electrical Power, p. 383 (2009).
120. Y. Hu, S. Yurkovich, Y. Guezennec and B. Yurkovich, *Control Engineering Practice*, 17, 1190 (2009).
121. Y. Hu, S. Yurkovich, Y. Guezennec and B. Yurkovich, *J Power Sources*, 196, 449 (2011).
122. L. Lam, P. Bauer and E. Kelder, in *2011 IEEE 33rd International Telecommunications Energy Conference (INTELEC)*, p. 1 (2011).
123. J. Zhang, S. Ci, H. Sharif and M. Alahmad, in *2010 Twenty-Fifth Annual IEEE Applied Power Electronics Conference and Exposition (APEC)*, p. 672 (2010).
124. B. Schweighofer, K. M. Raab and G. Brasseur, *IEEE transactions on instrumentation and measurement*, 52, 1087 (2003).
125. H. Chan, in *2000 IEEE Power Engineering Society Winter Meeting. Conference Proceedings (Cat. No. 00CH37077)*, p. 470 (2000).
126. T. Kim and W. Qiao, *IEEE Transactions on Energy Conversion*, 26, 1172 (2011).
127. T. Hu, B. Zanchi and J. Zhao, *IEEE Transactions on Energy Conversion*, 26, 787 (2011).
128. E. Kuhn, C. Forgez, P. Lagonotte and G. Friedrich, *J Power Sources*, 158, 1490 (2006).
129. E. Karden, S. Buller and R. W. De Doncker, *J Power Sources*, 85, 72 (2000).
130. S. Buller, M. Thele, R. W. De Doncker and E. Karden, *IEEE Transactions on Industry Applications*, 41, 742 (2005).
131. S. Buller, M. Thele, E. Karden and R. W. De Doncker, *J Power Sources*, 113, 422 (2003).
132. V. R. Subramanian, V. Boovaragavan and V. D. Diwakar, *Electrochem Solid St*, 10, A255 (2007).
133. V. R. Subramanian, V. Boovaragavan, V. Ramadesigan and M. Arabandi, *J Electrochem Soc*, 156, A260 (2009).
134. V. Ramadesigan, V. Boovaragavan, J. C. Pirkle Jr and V. R. Subramanian, *J Electrochem Soc*, 157, A854 (2010).
135. P. W. Northrop, V. Ramadesigan, S. De and V. R. Subramanian, *J Electrochem Soc*, 158, A1461 (2011).
136. V. Ramadesigan, K. Chen, N. A. Burns, V. Boovaragavan, R. D. Braatz and V. R. Subramanian, *J Electrochem Soc*, 158, A1048 (2011).
137. L. Cai and R. E. White, *J Electrochem Soc*, 156, A154 (2008).
138. M. Guo and R. E. White, *J Power Sources*, 198, 322 (2012).
139. B. Ng, P. T. Coman, W. E. Mustain and R. E. White, *J Power Sources*, 445, 227296 (2020).
140. Y. Dai, L. Cai and R. E. White, *J Electrochem Soc*, 160, A182 (2012).
141. Y. Dai, L. Cai and R. E. White, *J Electrochem Soc*, 161, E3348 (2014).
142. X. Jin, A. Vora, V. Hoshing, T. Saha, G. Shaver, R. E. García, O. Wasynczuk and S. Varigonda, *J Power Sources*, 342, 750 (2017).
143. Z. Chu, G. L. Plett, M. S. Trimboli and M. Ouyang, *Journal of Energy Storage*, 25, 100828 (2019).
144. Y. Yin, Y. Hu, S.-Y. Choe, H. Cho and W. T. Joe, *J Power Sources*, 423, 367 (2019).
145. L. Xia, E. Najafi, Z. Li, H. Bergveld and M. Donkers, *Appl Energ*, 208, 1285 (2017).
146. M. Farag, M. Fleckenstein and S. Habibi, *J Power Sources*, 342, 351 (2017).
147. X.-B. Meng, H.-X. Li and H.-D. Yang, *IEEE Transactions on Industrial Informatics*, 15, 2838 (2018).
148. A. Rodríguez, G. L. Plett and M. S. Trimboli, *eTransportation*, 1, 100009 (2019).
149. A. Farmann and D. U. Sauer, *Appl Energ.*, **225**, 1102 (2018).
150. A. B. Ansari, V. Esfahanian and F. Torabi, *Modares Mechanical Engineering*, **15**, 43 (2016).
151. Z. Li, D. Danilov, M. Donkers and H. J. Bergveld, in *Proc. 35th Benelux Meeting Syst. Control*, p. 43 (2016).
152. S. Yuan, L. Jiang, C. Yin, H. Wu and X. Zhang, *J Power Sources*, **352**, 245 (2017).
153. A. Subramaniam, S. Kolluri, C. D. Parke, M. Pathak, S. Santhanagopalan and V. R. Subramanian, *J Electrochem Soc*, **167**, 013534 (2019).
154. X. Zhao, Y. Yin, Y. Hu and S.-Y. Choe, *J Power Sources*, **418**, 61 (2019).
155. X. Zhao, Y. Bi and S.-Y. Choe, *SAE International Journal of Alternative Powertrains*, **6** (2017).
156. X. Zhao, Y. Bi, S.-Y. Choe and S.-Y. Kim, *Electrochim Acta*, **280**, 41 (2018).
157. X. Li, S.-Y. Choe and W. T. Joe, *J Power Sources*, **294**, 545 (2015).
158. Y. Zhao, S.-Y. Choe and J. Kee, *Electrochim Acta*, **270**, 440 (2018).
159. M. Song and S.-Y. Choe, *J Power Sources*, **436**, 226835 (2019).
160. Y. Yin and S.-Y. Choe, *Appl Energ.*, **271**, 115232 (2020).
161. Y. Bi, X. Zhao and S.-Y. Choe, in *2019 American Control Conference (ACC)*, p. 3155 (2019).
162. S. Alavi, C. Birkl and D. Howey, *J Power Sources*, **288**, 345 (2015).
163. G. H. Florentino, Model Predictive Control Implementations for Management of Lithium-ion Battery Cells and Packs Using Reduced-order Electrochemical Models, in, University of Colorado Colorado Springs. Kraemer Family Library (2018).
164. V. Ramadesigan, P. W. Northrop, S. De, C. M. Hoff, O. Leitermann, M. L. Crow, S. Santhanagopalan and V. R. Subramanian, *J Electrochem Soc*, **159**, R31 (2012).
165. Y. Cao, S. B. Lee, V. R. Subramanian and V. M. Zavala, *Journal of Process Control*, **90**, 46 (2020).
166. S. Kolluri, S. V. Aduru, M. Pathak, R. D. Braatz and V. R. Subramanian, *J Electrochem Soc*, **167**, 063505 (2020).
167. S. Dey and B. Ayalew, *Journal of Dynamic Systems, Measurement, and Control*, **139** (2017).
178. G. K. Singh, G. Ceder and M. Z. Bazant, *Electrochim Acta*, **53**, 7599 (2008).
179. W. Waag, S. Käbitz and D. U. Sauer, *Appl Energ*, **102**, 885 (2013).
180. M. Singh, J. Kaiser and H. Hahn, *Batteries*, **2**, 35 (2016).
181. J. Schmalstieg, S. Käbitz, M. Ecker and D. U. Sauer, *J Power Sources*, **257**, 325 (2014).
182. M. B. Pinson and M. Z. Bazant, *J Electrochem Soc*, **160**, A243 (2012).
183. M.-T. von Srbik, M. Marinescu, R. F. Martinez-Botas and G. J. Offer, *J Power Sources*, **325**, 171 (2016).
184. H. Lundgren, P. Svens, H. Ekström, C. Tengstedt, J. Lindström, M. Behm and G. Lindbergh, *J Electrochem Soc*, **163**, A309 (2015).
185. T. Waldmann, B.-I. Hogg, M. Kasper, S. Grolleau, C. G. Couceiro, K. Trad, B. P. Matadi and M. Wohlfahrt-Mehrens, *J Electrochem Soc*, **163**, A1232 (2016).
186. M. Mathew, S. Janhunen, M. Rashid, F. Long and M. Fowler, *Energies*, **11**, 1490 (2018).
187. K. S. Hariharan and V. S. Kumar, *J Power Sources*, **222**, 210 (2013).
188. S. Cho, H. Jeong, C. Han, S. Jin, J. H. Lim and J. Oh, *Computers & Chemical Engineering*, **41**, 1 (2012).
189. Y. Qi and S. J. Harris, *J Electrochem Soc*, **157**, A741 (2010).
190. S.-C. Yin, Y.-H. Rho, I. Swainson and L. Nazar, *Chem Mater*, **18**, 1901 (2006).
191. J. Shu, *J Solid State Elecr*, **13**, 1535 (2009).
192. F. Sagane, R. Shimokawa, H. Sano, H. Sakaëbe and Y. Iriyama, *J Power Sources*, **225**, 245 (2013).
193. L. Gireaud, S. Grugeon, S. Laruelle, B. Yrieix and J.-M. Tarascon, *Electrochim Commun*, **8**, 1639 (2006).
194. J. Steiger, D. Kramer and R. Mönig, *Electrochim Acta*, **136**, 529 (2014).
195. J. Guo, Z. Wen, M. Wu, J. Jin and Y. Liu, *Electrochem Commun*, **51**, 59 (2015).
196. K. J. Harry, D. T. Hallinan, D. Y. Parkinson, A. A. MacDowell and N. P. Balsara, *Nature materials*, **13**, 69 (2014).
197. A. Yermukhambetova, C. Tan, S. R. Daemi, Z. Bakenov, J. A. Darr, D. J. Brett and P. R. Shearing, *Scientific reports*, **6**, 35291 (2016).
198. F. Sun, L. Zielek, H. Markötter, A. Hilger, D. Zhou, R. Moroni, R. Zengerle, S. Thiele, J. Banhart and I. Manke, *ACS nano*, **10**, 7990 (2016).
199. D. Zhang, A. Forner-Cuenca, O. O. Taiwo, V. Yufit, F. R. Brushett, N. P. Brandon, S. Gu and Q. Cai, *J Power Sources*, **447**, 227249 (2020).
200. R. F. Ziesche, T. Arlt, D. P. Finegan, T. M. Heenan, A. Tengattini, D. Baum, N. Kardjilov, H. Markötter, I. Manke and W. Kockelmann, *Nat Commun*, **11**, 1 (2020).
201. B. Key, R. Bhattacharyya, M. Morcrette, V. Seznec, J.-M. Tarascon and C. P. Grey, *Journal of the American Chemical Society*, **131**, 9239 (2009).
202. F. Chevallier, M. Letellier, M. Morcrette, J.-M. Tarascon, E. Frackowiak, J.-N. Rouzaud and F. Béguin, *Electrochim Solid St*, **6**, A225 (2003).
203. M. Letellier, F. Chevallier and M. Morcrette, *Carbon*, **45**, 1025 (2007).
204. R. Gerald Ii, C. Johnson, J. Rathke, R. Klingler, G. Sandi and L. Scanlon, *J Power Sources*, **89**, 237 (2000).
205. Y.-C. Hsieh, M. Leißing, S. Nowak, B.-J. Hwang, M. Winter and G. Brunklaus, *Cell Reports Physical Science*, 100139 (2020).
206. M. Leskes, G. Kim, T. Liu, A. L. Michan, F. Aussenac, P. Dorffer, S. Paul and C. P. Grey, *The Journal of Physical Chemistry Letters*, **8**, 1078 (2017).
207. Y. Jin, N.-J. H. Kneusels, P. C. Magusin, G. Kim, E. Castillo-Martínez, L. E. Marbella, R. N. Kerber, D. J. Howe, S. Paul and T. Liu, *Journal of the American Chemical Society*, **139**, 14992 (2017).

208. Y. Jin, N.-J. H. Kneusels, L. E. Marbella, E. Castillo-Martínez, P. C. Magusin, R. S. Weatherup, E. Jónsson, T. Liu, S. Paul and C. P. Grey, *Journal of the American Chemical Society*, **140**, 9854 (2018).

209. M. A. Hope, B. L. Rinkel, A. B. Gunnarsdóttir, K. Märker, S. Menkin, S. Paul, I. V. Sergeyev and C. P. Grey, *Nat Commun*, **11**, 1 (2020).

210. F. Rosciano, M. Holzapfel, W. Scheifele and P. Novák, *Journal of Applied Crystallography*, **41**, 690 (2008).

211. M. Roberts, J. J. Biendicho, S. Hull, P. Beran, T. Gustafsson, G. Svensson and K. Edström, *J Power Sources*, **226**, 249 (2013).

212. T. Kulova, A. Skundin, E. Nizhnikovskii and A. Fesenko, *Russian Journal of Electrochemistry*, **42**, 259 (2006).

213. G. Feher and A. Kip, *Physical Review*, **98**, 337 (1955).

214. J. H. Pifer and R. Magno, *Physical Review B*, **3**, 663 (1971).

215. M. Sathiya, J.-B. Leriche, E. Salager, D. Gourier, J.-M. Tarascon and H. Vezin, *Nat Commun*, **6**, 1 (2015).

216. A. Niemöller, P. Jakes, R.-A. Eichel and J. Granwehr, *Scientific reports*, **8**, 1 (2018).

217. F. J. Dyson, *Physical Review*, **98**, 349 (1955).

218. P. Glover and P. Mansfield, *Reports on progress in physics*, **65**, 1489 (2002).

219. M. Weiger, D. Schmidig, S. Denoth, C. Massin, F. Vincent, M. Schenkkel and M. Fey, *Concepts in Magnetic Resonance Part B: Magnetic Resonance Engineering: An Educational Journal*, **33**, 84 (2008).

220. A. Webb, *Concepts in Magnetic Resonance Part A: An Educational Journal*, **22**, 25 (2004).

221. G. Maresch, M. Mehring and S. Emid, *Physica B+ C*, **138**, 261 (1986).

222. M.-T. F. Rodrigues, G. Babu, H. Gullapalli, K. Kalaga, F. N. Sayed, K. Kato, J. Joyner and P. M. Ajayan, *nature energy*, **2**, 1 (2017).

223. A. Mistry, C. Fear, R. Carter, C. T. Love and P. P. Mukherjee, *ACS Energy Letters*, **4**, 156 (2018).

224. R. Carter, B. Huhman, C. T. Love and I. V. Zenyuk, *J Power Sources*, **381**, 46 (2018).

225. F. B. Spingler, W. Wittmann, J. Sturm, B. Rieger and A. Jossen, *J Power Sources*, **393**, 152 (2018).

226. Q. Liu, D. Xiong, R. Petibon, C. Du and J. Dahn, *J Electrochem Soc*, **163**, A3010 (2016).

227. A. Hsieh, S. Bhadra, B. Hertzberg, P. Gjeltema, A. Goy, J. W. Fleischer and D. A. Steingart, *Energy & environmental science*, **8**, 1569 (2015).

228. G. Davies, K. W. Knehr, B. Van Tassell, T. Hodson, S. Biswas, A. G. Hsieh and D. A. Steingart, *J Electrochem Soc*, **164**, A2746 (2017).

229. L. Göld, T. Bach, W. Virsik, A. Schmitt, J. Müller, T. E. Staab and G. Sextl, *J Power Sources*, **343**, 536 (2017).

230. J. B. Robinson, M. Maier, G. Alster, T. Compton, D. J. Brett and P. R. Shearing, *Physical Chemistry Chemical Physics*, **21**, 6354 (2019).

231. Z. Deng, Z. Huang, Y. Shen, Y. Huang, H. Ding, A. Luscombe, M. Johnson and J. Dahn.
232. C. Bommier, W. Chang, Y. Lu, J. Yeung, G. Davies, R. Mohr, M. Williams and D. Steingart, *Cell Reports Physical Science*, 100035 (2020).
233. W. Chang, R. Mohr, A. Kim, A. Raj, G. Davies, K. Denner, J. H. Park and D. Steingart, *Journal of Materials Chemistry A*, 8, 16624 (2020).
234. S. E. Renfrew and B. D. McCloskey, *Journal of the American Chemical Society*, 139, 17853 (2017).
235. B. D. McCloskey, A. Valery, A. C. Luntz, S. R. Gowda, G. M. Wallraff, J. M. Garcia, T. Mori and L. E. Krupp, *The journal of physical chemistry letters*, 4, 2989 (2013).
236. S. E. Renfrew and B. D. McCloskey, *ACS Applied Energy Materials*, 2, 3762 (2019).
237. M. M. Ottakam Thotiyl, S. A. Freunberger, Z. Peng and P. G. Bruce, *Journal of the American Chemical Society*, 135, 494 (2013).
238. X.-G. Yang, G. Zhang, S. Ge and C.-Y. Wang, *Proceedings of the National Academy of Sciences*, 115, 7266 (2018).