Research Article

Enhanced Light Emission from Type-II Red InGaN/GaNSb/GaN Quantum-Well Structures

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Received 23 May 2022; Accepted 11 August 2022; Published 26 September 2022

Academic Editor: Da-Ren Hang

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Electronic and optical properties of type-II InGaN/GaNSb/GaN quantum-well (QW) structures are investigated by using the multiband effective mass theory for potential applications in red light-emitting diodes. The heavy-hole effective mass around the topmost valence band is not affected much by the insertion of the GaNSb layer, and the optical matrix elements are greatly increased by the inclusion of the GaNSb layer in the InGaN/GaN QW structure. As a result, the type-II InGaN/GaNSb/GaN QW structure shows a much larger emission peak than the conventional type-I QW structure owing to the decrease in spatial separation between electron and hole wavefunctions, in addition to the reduction of the effective well width. It is also observed that the In content in InGaN well can be significantly reduced for the type-II QW structure with a large Sb content, compared to that for the type-I QW structure.

1. Introduction

III-nitride-based white light-emitting diodes (LEDs) have received much attention for their tremendous potential for energy-efficient general illumination applications [1–3]. They have been usually fabricated by using InGaN/GaN quantum wells (QWs) as InGaN has a direct band gap covering the entire visible spectrum. So far, blue and green InGaN/GaN LEDs have been improved greatly although the efficiency of green LEDs is still inferior to that of blue LEDs. On the other hand, emission in the red wavelength range typically exhibits much lower quantum efficiency because higher In content and relatively thicker QWs are necessary. [4] The high In content in the QW region leads to the large piezoelectric (PZ) and spontaneous (SP) polarization fields [5, 6], which significantly reduce the internal quantum efficiency of the QW owing to a larger spatial separation between electron and hole wavefunctions.

As a method to reduce the internal field effect, type-II QW structures have been studied by several research groups. The structures considered include the type-II InGaN/GaNAS QW [7, 8], the type-II InGaN/ZnGeN₂ QW [9], and the type-II InGaN/ZnSnN₂/GaN QW [10, 11]. In addition, the type-II InGaN/GaNSb/GaN QW was proposed to improve the internal quantum efficiency of the InGaN/GaN QW structure for the blue wavelength range [12]. On the other hand, in the case of the red wavelength range, there has been no work reported on the optical properties of the type-II QW structures combining InGaN and GaNSb.

In this paper, we theoretically investigate the electronic and optical properties of the type-II QW structure combining In₁ₓGa₁₋ₓN and GaN₁₋ᵧSbᵧ as one way of improving the internal quantum efficiency of the red InₓGa₁₋ₓN/GaN QW structure. The multiband effective mass theory is utilized to calculate the optical properties of these InGaN/GaNSb/GaN QW structures. The finite difference method (FDM) is implemented in Fortran codes to calculate the valence-band structures. These results are compared with those of the conventional red InGaN/GaN QW structure. The self-consistent (SC) band structures and wavefunctions...
2. Results and Discussion

Figure 1 shows potential energy profiles for (a) type-I In\(_{x}\)Ga\(_{1-x}\)N/GaN (\(x = 0.47\)) and (b) type-II In\(_{x}\)Ga\(_{1-x}\)N/GaSb\(_{y}\)/GaN (\(x = 0.45, y = 0.04\)) QW structures. The wavefunctions of the first conduction subband (C1) and the first valence subband (HH1) at the zone center are also shown as blue and red lines, respectively.

Figure 1: Potential energy profiles for (a) type-I In\(_{x}\)Ga\(_{1-x}\)N/GaN (\(x = 0.47\)) and (b) type-II In\(_{x}\)Ga\(_{1-x}\)N/GaSb\(_{y}\)/GaN (\(x = 0.45, y = 0.04\)) QW structures. The wavefunctions of the first conduction subband (C1) and the first valence subband (HH1) at the zone center are also shown as blue and red lines, respectively.

are obtained by solving the Schrödinger equation for electrons, the 3 × 3 Hamiltonian for holes, and the Poisson’s equation iteratively [13, 14]. The non-Markovian free-carrier model with bandgap renormalization to calculate the spontaneous emission spectrum \(g_{sp}(\omega)\) is taken from [15] and [16]. The bandgap energy of GaN\(_{1-x}\)Sb\(_{y}\) used in the calculation is obtained by \(E_{GaSb}^g(y) = yE_{GaSb}^g(1 - y)E_{GaN}^g - Cy(1 - y)\), where \(E_{GaSb}^g\) and \(E_{GaN}^g\) are bandgap energies of GaSb and GaN, respectively, and \(C\) is the bowing parameter. Here, we assume that the unknown parameters of GaNSb are same as those of GaN as a first approximation. The material parameters for GaN and InN used in the calculation are taken from [17] and references therein.

In the case of the conventional type-I QW structure, a large spatial separation between electron and hole wavefunctions is observed due to the large internal field. On the other hand, a spatial separation between electron and hole wavefunctions in the type-II InGaN/GaNSb/GaN QW structure is greatly reduced, compared to that of type-I QW structure. Thus, we expect that the optical transition probability will be greatly improved with the inclusion of the GaNSb layer, as discussed in the following.

Figure 2 shows the valence-band structures where Figure 2(a) is for the type-I InGaN/GaN QW and Figure 2(b) is for the type-II InGaN/GaNSb/GaN QW, considered in Figure 1. The SC solutions are obtained at a sheet carrier density of \(N_{2D} = 20 \times 10^{12}\) cm\(^{-2}\). The heavy-hole effective mass around the topmost valence band is not affected much by the GaNSb layer. The overall shapes of the valence-band structures for both cases are similar to each other. However, the energy spacing between the first two subbands and the third two subbands of the type-II QW structure is slightly smaller than that of the type-I QW structure.

Figure 3 shows transverse-electric (TE)-polarized optical matrix elements where Figure 3(a) is for the type-I InGaN/GaN QW and Figure 3(b) is for the type-II InGaN/GaNSb/GaN QW as a function of \(k_y\) at sheet carrier densities of 2 and \(20 \times 10^{12}\) cm\(^{-2}\). Here, we use 2 and \(20 \times 10^{12}\) cm\(^{-2}\) as examples of relatively low and high sheet carrier densities. TE-polarized optical matrix elements are shown to be nearly independent of the in-plane wave vector \(k_y\) for both QW structures. In the case of the low sheet carrier density, the
Figure 2: Valence-band structures for (a) type-I \( \text{In}_{x}\text{Ga}_{1-x}\text{N}/\text{GaN} \) \( x = 0.47 \) and (b) type-II \( \text{In}_{x}\text{Ga}_{1-x}\text{N}/\text{GaSbN}/\text{GaN} \) \( x = 0.45, \text{Sb} = 0.04 \) QW structures.

Figure 3: TE-polarized optical matrix elements for (a) type-I \( \text{In}_{x}\text{Ga}_{1-x}\text{N}/\text{GaN} \) \( x = 0.47 \) and (b) type-II \( \text{In}_{x}\text{Ga}_{1-x}\text{N}/\text{GaSbN}/\text{GaN} \) \( x = 0.45, \text{Sb} = 0.04 \) QW structures as a function of \( k_{||} \) at sheet carrier densities of 1\( \times \) and 20\( \times \)10\(^{12}\) cm\(^{-2}\).
matrix element of the type-II InGaN/GaNSb/GaN QW structure is shown to be smaller than that of the conventional type-I QW structure. However, at the higher sheet carrier density, the type-II InGaN/GaNSb/GaN QW structure shows a larger matrix element than the conventional type-I QW structure because the reduction effect in the built-in potential by the carrier screening is dominant for the type-II QW structure.

Figure 4 shows TE-polarized spontaneous emission spectra (Figure 4(a)) and peak intensities as a function of sheet carrier density (Figure 4(b)) for type-I InGaN/GaN and type-II In$_{x}$Ga$_{1-x}$/N/GaSb$_{y}$/N/GaN ($x=0.45$, $y=0.04$) QW structures. TE-polarized spontaneous emission spectra are obtained at a sheet carrier density of $20 \times 10^{12}$ cm$^{-2}$. The peak wavelengths of both QW structures are shown to be approximately 640 nm. The type-II InGaN/GaNSb/GaN QW structure shows a larger emission peak than the conventional type-I QW structure in the investigated range of sheet carrier densities although the matrix element of the type-II InGaN/GaNSb/GaN QW structure is smaller than that of the conventional type-I QW structure owing to a small screening effect. This is mainly because the density of states $k_f/\left(\pi m^*_{eff}\right)$ for the type-II QW structure is enhanced owing to the reduction of the effective well width, $m^*_{eff}$. The effective well width is given by $m^*_{eff} = \left[m^*_{w2} (L_{w2} + L_{v2})/2 + m^*_{w1} L_{w2}\right]/(m^*_{w2} + m^*_{w1})$, where $m^*_{w1}$ and $m^*_{w2}$ are hole and electron effective masses for wells, respectively [18].

In particular, in the range of higher sheet carrier densities, the peak intensity of the type-II InGaN/GaNSb/GaN QW structure is much larger than that of the conventional type-I QW structure owing to an increased matrix element. Also, the peak intensities of both QW structures rapidly increase with increasing sheet carrier density. The increasing rate of the peak intensity for type-II QW structure is larger than that for type-I QW structure. If the Sb content in GaNSb is increased, the In content in InGaN well can be significantly reduced, as discussed in the following.

Figure 5 shows the potential energy profile for the type-II In$_{x}$Ga$_{1-x}$/N/GaSb$_{y}$/N/GaN QW structure with $x=0.47$ (Figure 5(a)) and TE-polarized spontaneous emission spectra of both type-I and type-II QW structures (Figure 5(b)). The inset in Figure 5(b) shows the TE-polarized optical matrix element for the type-II In$_{x}$Ga$_{1-x}$/N/GaSb$_{y}$/N/GaN QW structure ($x=0.3$, $y=0.08$). The SC solutions are obtained at a sheet carrier density of $20 \times 10^{12}$ cm$^{-2}$. The thickness of the GaNSb layer is set to be 1 nm as in Figure 1, but the Sb content of the GaNSb layer is changed from 0.04 to 0.08. With this change, the transition wavelength of 640 nm is obtained at an In content of 0.3 in the InGaN well. This value is much smaller than that for the type-I In$_{x}$Ga$_{1-x}$/N/GaN QW structure ($x=0.47$). In addition, we observe that the type-II InGaN/GaNSb/GaN QW structure still has a much larger emission peak than the conventional type-I QW structure even for the case of...
smaller In content. On the other hand, the emission peak of the type-II $\text{In}_x\text{Ga}_{1-x}\text{N}/\text{GaSb}_y\text{N}/\text{GaN}$ structure ($x=0.45$, $y=0.04$) is slightly smaller than that of the type-II $\text{In}_x\text{Ga}_{1-x}\text{N}/\text{GaSb}_y\text{N}/\text{GaN}$ structure ($x=0.3$, $y=0.08$) because the former has a smaller optical matrix element than the latter owing to the reduced overlap between electron and hole wavefunctions.

3. Summary

In summary, electronic and optical properties of type-II red InGaN/GaNSb/GaN QW structure have been investigated by using the multiband effective mass theory. Optical matrix elements for the type-II InGaN/GaNSb/GaN QW structures are greatly increased owing to the screening effect with increasing sheet carrier density. The type-II InGaN/GaNSb/GaN QW structures show much larger emission peaks than the conventional type-I QW structure in an investigated range of sheet carrier densities because the density of states for the type-II QW structure is enhanced owing to the reduction of the effective well width. Also, we have observed that the In content in InGaN well can be significantly reduced with increasing Sb content in the GaNSb layer.

Data Availability

The data used to support the findings of this study are available from the corresponding author upon request.

Conflicts of Interest

The authors declare that they have no conflicts of interest.

Acknowledgments

This research was supported by the Nano Material Technology Development Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Science and ICT (Grant 2021M3D1A2048623).

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