HE 1113-0641: THE SMALLEST-SEPARATION QUADRUPE LENS IDENTIFIED BY A GROUND-BASED OPTICAL TELESCOPE*†

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ABSTRACT

The Hamburg/ESO quasar HE 1113-0641 is found to be a quadruple gravitational lens, based on observations with the twin 6.5 m Magellan telescopes at the Las Campanas Observatory, and subsequently with the Hubble Space Telescope. The \( z_s = 1.235 \) quasar appears in a cross configuration, with \( i' \) band magnitudes ranging from 18.0 to 18.8. With a maximum image separation of 0.67, this is the smallest-separation quadruple ever identified using a ground-based optical telescope. Point-spread function (PSF) subtraction reveals a faint lensing galaxy. A simple lens model succeeds in predicting the observed positions of the components, but fails to match their observed flux ratios by up to a magnitude. We estimate the redshift of the lensing galaxy to be \( z_L \sim 0.7 \). Time delay estimates are on the order of a day, suggesting that the flux ratio anomalies are not due to variability of the quasar, but may result from substructure or microlensing in the lens galaxy.

Key words: gravitational lensing – quasars: individual (HE 1113-0641)

1. INTRODUCTION

A quasar lensed by a galaxy along the line of sight presents opportunities to study both the gravitational potential that is doing the lensing and the quasar that is lensed. In particular the inevitable microlensing of the quasar by stars in the lensing galaxies permits estimates of the stellar (as opposed to dark) matter surface density of the lensing galaxy (Schechter & Wambsganss 2004) and of the size of the quasar accretion disk (Pooley et al. 2007; Poindexter et al. 2007) and emission line region (Keeton et al. 2006). Quadruply imaged systems are especially useful for such purposes, giving multiple lines of sight through the lens and additional constraints on the lensing potential.

But the rate at which new systems are discovered has leveled off to something like a half dozen per year, with the most recent discoveries coming largely from the Sloan Digital Sky Survey (SDSS) (e.g., Oguri et al. 2007). Their selection algorithm is incomplete at separations below 1 arcsec and worse for quadruple systems than for doubles (Oguri et al. 2006). Of the new SDSS lenses, just three are quadruple. No comparable survey has been carried out in the southern hemisphere.

We report here the discovery of a new quadruply imaged system, HE 1113-0641. It was observed as part of a program to identify lensed systems among the Hamburg/ESO quasar survey (Wisotzki et al. 2000). While included within the SDSS survey area, it has not, until the present work, been identified as a lensed system.

In Section 2, we report the observations made using Magellan and the Hubble Space Telescope (HST). Section 3 describes our analysis of the data. In Section 4 we describe a simple model of the lens, and in Section 5 we make a rough estimate of the lens redshift. In Section 6 we discuss the conclusions we can come to regarding HE 1113-0641. Throughout this paper, we adopt a "concordance" cosmology with \( \Omega_M = 0.3, \Omega_\Lambda = 0.7 \), and \( h = 0.7 \).

2. OBSERVATIONS

HE 1113-0641 was originally discovered to be a \( z = 1.235 \) quasar in the Hamburg/ESO digital objective prism survey (Wisotzki et al. 2000). Based on its redshift and apparent magnitude \( B = 17.01 \), it was found to have a relatively high lensing probability, and was selected for a follow-up observation.

We observed the object in early 2002 and early 2003 in the Sloan \( g' \) and \( i' \) bands using the Magellan 6.5 m telescopes. In fall 2003, we observed it in three bands using the Advanced Camera for Surveys (ACS) and Near-Infrared Camera and Multi-Object Spectrometer (NICMOS) detectors on the HST. These observations are tabulated in Table 1.

### Table 1

| Date         | Instrument | Filter | Exposure | Seeing |
|--------------|------------|--------|----------|--------|
| 2002 February 16 | MagIC | \( g' \) | 7 × 60 s | 0".52 |
| 2003 January 26   | MagIC | \( i' \) | 7 × 60 s | 0".43 |
| 2003 November 06  | NICMOS | \( H \) | 3 × 640 s | ⋯ |
| 2003 November 07  | ACS/WFC | \( V \) | 3 × 120 s | ⋯ |
| 2003 November 07  | ACS/WFC | \( I \) | 2 × 480 s | ⋯ |

2.1. Magellan 6.5 m Imaging

HE 1113-0641 was identified as a quadruple gravitational lens on 2002 February 16 using the Baade 6.5 m telescope at the Las Campanas Observatory. Seven 60 s exposures in each of the Sloan \( i' \) and \( g' \) bands, and a single exposure in the \( r' \) band, were taken using the Raymond and Beverly Sackler Magellan
2.2. Hubble Space Telescope Imaging

On 2003 November 6 and 7, HE 1113-0641 was observed using both the NIC2 camera of the NICMOS and the Wide Field channel (WFC) of the ACS. The NICMOS images had 256 × 256 pixels and a 19′2 field of view, while those from the ACS were significantly larger, with 4096 × 4096 pixels filling a 3.4 arcmin field of view. Three filters were used, F160W with NICMOS and F555W and F814W with the ACS (hereafter \( H \), \( V \), and \( I \), respectively). Because of the diffraction-limited quality of the images, they were not well sampled, with the width of PSF ranging from 1.5 pixels in \( H \) to 2.1 pixels in \( I \).

We used the Multidrizzle program of Koekemoer et al. (2002b), version 2.2, to register the ACS images, clean them of cosmic rays, and drizzle them onto a single image per filter. The drizzling process also corrects for geometric distortion arising from the design of the camera. We likewise drizzled the NICMOS images into a single image using the procedure detailed in Koekemoer et al. (2002a).3

The drizzled ACS and NICMOS images of HE 1113-0641 may be seen in Figure 2.

3. ANALYSIS

The small separation of this lens, combined with the relative faintness of the lensing galaxy, complicated the task of disentangling the four quasar components and the galaxy, particularly for the ground-based data. To address this issue, we fitted for the relative positions and brightnesses of the objects using an iterative process. First we performed a fit to each image for the relative positions of the objects, then averaged the values and repeated the fit, holding constant the relative positions, to determine the photometry.

We used a variant of the DoPHOT photometry package (Schechter et al. 1993) called Clumpfit to carry out the fits, using empirical PSFs provided by field stars for the quasar components and a pseudo-Gaussian profile for the lensing galaxy. (Though this is not a physical profile choice, we found that the choice

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3 The HST Dither Handbook (Koekemoer et al. 2002a) is available at http://www.stsci.edu/hst/wfpc2/documents/dither_handbook.html.
of galaxy profile had a negligible effect on the goodness of the fit.) We also used DoPHOT to obtain astrometry and aperture photometry for several other stars in the wider-field (ACS and Magellan) images.

3.1. Magellan data

A fit consisting of four empirical PSFs (provided by a field star) was made to the stacked $i'$ band Magellan images. This came to a total of 13 free parameters: two-dimensional position and normalization for all four objects, and the sky level. It was clear from the residual images that a small amount of leftover flux remained; indeed, that the four point sources had been over-subtracted in an attempt to compensate (see Figure 3). So a second fit was attempted using a model with four point sources and a circularly symmetric galaxy; however, there was not enough galaxy flux for the fit to distinguish between radial profiles or converge upon a scale size. We therefore chose a fixed-width pseudo-Gaussian profile for the galaxy and repeated the fit, which now had 16 free parameters. The width of the galaxy was set to the width chosen for the ACS data (see Section 3.2.1), suitably broadened by the PSF. This fit was able to account for the leftover flux.

The relative astrometry resulting from the 16-parameter fits in $i'$ and other filters is listed in Table 2. The relative positions of the four quasar components were then weighted by the inverse of their uncertainties and averaged, yielding final values, which may be found in Table 5. In the case of the lensing galaxy, only the HST positions were averaged (see Section 3.2.1).

Once the relative astrometry had thus been determined, we repeated the fit with fixed relative positions and a fixed galaxy width. The results of this fit may be seen in Table 4, and residual images are in Figure 2.

For the $g'$ band images, which had poorer seeing, the quasar components were too blurred for a successful fit until relative positions were fixed. The residual images of these fits may also be seen in Figure 2. There was no indication in the residuals of leftover flux indicative of a lens galaxy, so we conclude that we have not detected it in the $g'$ band.

Aperture photometry was also obtained for several field stars, including those used as model PSFs. The positions of these stars, labeled a through e in Figure 1, may be found in Table 3, and their magnitudes are listed in Table 4.

To enable absolute flux calibration, aperture photometry was obtained for standard stars from the sample of Smith et al. (2002). For the 2002 data we used PG 1047+003A, and for that of 2003 we used RU 152. We applied a first-order correction for atmospheric extinction when calculating the zero points, using extinction coefficients from Table 4 of Smith et al. (2007). It is worth noting that the $u'g'r'i'z'$ system is a broadband approximation to the (monochromatic) AB magnitude system, and is given by

$$m = -2.5 \log \frac{\int d(\log v) f_\nu S_\nu}{\int d(\log v) S_\nu} - 48.60,$$

which is calibrated by synthetic spectra of BD+17 4708. The $u'g'r'i'z'$ system deviates from the true AB system by less than 5% (Smith et al. 2002), and is presented in Table 4.

We estimate the uncertainty in the relative photometry to be 0.1 magnitudes in $g'$ and 0.05 magnitudes in $i'$. Absolute photometry is less certain, with error bars a factor of 1.4 larger. With these uncertainties the data are consistent with a slight overall brightening of all four images between 2002 and 2003, probably caused by intrinsic variability of the quasar. However, they fail to convincingly demonstrate uncorrelated changes in the flux ratios over time, even when combined with HST data; such variations might have been indicative of microlensing.

Despite our use of fixed positions for the $g'$ band images, there are inconsistencies in the $g' - i'$ colors of the quasar components. It is likely that these are due to the confusion caused by mediocre seeing in the $g'$ band.

There can be little doubt that the lensing galaxy has been detected in the $i'$ band in both data sets. However, its size and shape remain poorly constrained. By fitting a fixed pseudo-Gaussian to both epochs of data, we were able to estimate its $i'$ band flux, but with substantial uncertainty (0.2 magnitudes

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### Table 2

| B         | C         | D         | G         |
|-----------|-----------|-----------|-----------|
| $\Delta \alpha \cos (\delta)$ | $\Delta \delta$ | $\Delta \alpha \cos (\delta)$ | $\Delta \delta$ |
| $i'$ (2002) | $-0.515$ $+0.428$ | $-0.148$ $+0.433$ | $-0.431$ $+0.188$ |
| $i'$ (2003) | $-0.517$ $+0.424$ | $-0.149$ $+0.432$ | $-0.422$ $+0.134$ |
| $V$       | $-0.518$ $+0.424$ | $-0.152$ $+0.427$ | $\cdots$ $\cdots$ |
| $I$       | $-0.519$ $+0.422$ | $-0.152$ $+0.429$ | $-0.320$ $+0.145$ |
| $H$       | $-0.518$ $+0.425$ | $-0.146$ $+0.429$ | $-0.308$ $+0.169$ |

**Note.** All positions are given relative to component A.

### Table 3

| Field Stars a through e | $\Delta \alpha \cos (\delta)$ | $\Delta \delta$ |
|------------------------|--------------------------|------------------|
| a                      | $+0.04$                  | $-66.50$        |
| b                      | $-5.70$                  | $-79.16$        |
| c                      | $+66.47$                 | $-29.91$        |
| d                      | $+85.97$                 | $-29.07$        |
| e                      | $+44.84$                 | $+39.90$        |

**Note.** All positions are given relative to component A.
of difference between epochs). We were able to determine the position of the galaxy using HST data, but its size and shape remained elusive (see Section 3.2).

3.2. HST Data

The HST images did not suffer from inadequate seeing, but rather from under-sampling of the PSF, leading to complications in the interpolation of empirical PSFs. We therefore resampled the ACS images to a scale of 0.03 per pixel, and the NICMOS image to a scale of 0.0375 per pixel, when drizzling.

3.2.1. ACS

Since the ACS PSF is known to vary across the field of view and also with time, we searched the HST archive for images with a suitable PSF star located close to the position of HE 1113-0641 on the chip, and obtained at a time close to 2003 November 7. In the V band, we used a field obtained on 2003 October 7, and in I we used a field obtained on 2003 November 25. In order to minimize differences in the PSF caused by the drizzling process, we used the same Multidrizzle process on these images as on the HE 1113-0641 images. We chose PSF stars that were close to the correct position on the chip, and were not saturated. In both cases, these stars were about 1 magnitude fainter than image A.

The fits proceeded as they had in the case of the Magellan images. There was appreciable leftover flux in the I band, concentrated in the center of the lens system (see Figure 3). Since the noisy residuals of the quasar components again prevented a measurement of the lensing galaxy’s radial profile or scale size, we fit it as a circular pseudo-Gaussian with a fixed width of 0.35 (broadened slightly by the ACS PSF). The width was chosen by the inspection of the residual image, since the choice had little to no effect on the goodness of the fit parameter. No sign of the galaxy was visible in the V band residual image.

Both of these fits were repeated once we had determined and fixed the relative positions. The resultant residual images are visible in Figure 2. The magnitudes were calibrated using AB zero-point keywords from the HST data headers. The HST broadband flux calibration is based on synthetic spectra of four primary white dwarf stars (Bohlin 1995), and agrees with the AB zero points of Smith et al. (2002) to within 3% (Bohlin & Gilliland 2004). The photometric data are presented in Table 4.

Astrometric measurements were made using DoPHOT on the I band image. A plate solution was found using 16 USNO-B stars. This solution gives the position of component A as (11 h 16m 23.56, −06° 57′ 38.6′′; J2000) to a precision of 0.01 in α and 0.1 in δ.

3.2.2. NICMOS

The analysis of the NICMOS data was similar to that of the ACS data. A 2003 September 5 observation of SA 107-626 provided a model PSF. In this case, the PSF star was much brighter than the quasar components.

Leftover flux from the lensing galaxy was visible in the H band image, as it had been in the I band image (see Figure 3). We again modeled it as a pseudo-Gaussian, with the same fixed width (broadened slightly by the NICMOS PSF).

A second fit was performed with the relative positions of the quasar components fixed to the averaged value. The photometry that resulted from this second fit was calibrated to the AB system using keywords from the data headers, and is visible in Table 4. The residual image may be seen in Figure 2.

4. MODELING THE LENS

Using the Lensmodel software of Keeton (2001), we modeled the lensing galaxy as a singular isothermal sphere plus external shear. This model consists of a projected two-dimensional lensing potential given by

$$\Psi_{2D}(\theta) = br - \frac{\gamma}{2} r^2 \cos(2(\phi - \phi_r)),$$

where $b$ is the monopole Einstein radius in arcseconds, $r$ and $\phi$ are the radial and angular components on the sky of the position vector $\theta$, and $\gamma$ and $\phi_r$ are the strength and direction of the external shear. Note that in this convention the position of a companion mass causing a shear would be along the $\pm \phi_r$ direction. (No such perturber is observed in this case, consistent with the apparently small shear strength.) This model has seven free parameters, and was constrained by the averaged positions for the four components and the galaxy.

We found the best model to have a radius of 0′′.332, with $\gamma = 0.04$ and $\phi_r = 37.7$ degrees east of north. The source

| g′ (2002) | g′ (2003) | i′ (2002) | i′ (2003) | V | I | H | g′ − i′ (2003) | V − I | I − H |
|---------|----------|----------|----------|---|---|---|----------------|-------|-------|
| A       | 18.37    | 18.19    | 18.02    | 17.96   | 18.33 | 18.32 | 18.25       | −0.23   | +0.00  | +0.07  |
| B       | 18.28    | 18.24    | 18.09    | 18.02   | 18.40 | 18.35 | 18.27       | −0.22   | +0.05  | +0.08  |
| C       | 18.53    | 18.39    | 18.46    | 18.37   | 18.64 | 18.61 | 18.74       | +0.02   | +0.03  | −0.13  |
| D       | 18.91    | 18.91    | 18.79    | 18.76   | 19.06 | 19.01 | 18.92       | +0.15   | +0.06  | +0.08  |
| G       | …        | …        | 22.36    | 22.17   | …    | …    | 22.47       | 21.05   | …     | 1.42   |
| a       | 20.63    | 20.70    | 18.08    | 18.05   | 20.16 | 17.86 | …          | +2.65   | +2.30  | …     |
| b       | …        | …        | 18.39    | 20.66   | 18.04 | …    | …          | +2.62   | …     | …     |
| c       | 19.29    | 19.34    | 16.62    | 16.62   | 18.91 | 16.39 | …          | +2.72   | +2.52  | …     |
| d       | 18.78    | 18.77    | 18.36    | 18.36   | 18.69 | 18.45 | …          | +0.41   | +0.24  | …     |
| e       | 22.51    | 22.59    | 20.01    | 20.02   | 22.06 | 19.81 | …          | +2.57   | +2.25  | …     |

Notes.

a All magnitudes are in the AB photometric system (see Section 3.1).

b Labels a through e are field stars.
position was predicted to be \((\Delta \alpha \cos(\delta), \Delta \delta) = (-0'.308, 0'.151)\) relative to the position of component A. Table 5 contains a summary of the model’s predictions, compared to observed data.

The model fits the component positions very well, even with the tight constraints provided by the \textit{HST}. It does not, however, correctly predict the flux ratios. This is in keeping with experience; optical flux anomalies can be seen in many lensed quasars (Keeton et al. 2003). In particular, the D component, a saddle point image, is predicted to be the brightest, but is observed to be the faintest, too faint by as much as a factor of 2.5. The predicted position of the lensing galaxy is 0'01 from the observed position. This is well within the estimated measurement error.

Finally, the model allows us to predict the time delays between the images, given a lens redshift. We used \(z_L = 0.7\), as estimated in Section 5. This yields the predicted time delays seen in Table 5, with the maximum delay being \(\sim 1.5\) days. We also calculated predicted time delays for \(z_L = 0.4\) and \(z_L = 1.0\); these changes reduced and increased (respectively) the time delays by a factor of \(\sim 3\). This strong dependence suggests that a measurement of the time delays might constrain the lens redshift; however, the unknown radial mass profile of the lensing galaxy is likely to have a similarly strong effect on time delays.

By way of comparison, HE 0435-1223, which has a shape similar to that of HE 1113-0641 but a larger image separation, has a maximum time delay of two weeks (Kochanek et al. 2006).

### 5. ESTIMATING THE LENS REDSHIFT

In order to estimate the redshift of the lensing galaxy and determine what its \(I - H\) color could tell us about its morphology, we combined the results of our lens model with properties of typical galaxies.

From the lens strength \(b = 0'.332\) we found the line-of-sight velocity dispersion of the lens using

\[
b = \frac{D_{LS}}{D_S} \frac{4\pi\sigma^2}{c^2}
\]

(Narayan & Bartelmann 1996), where \(D_{LS}\) is the angular diameter distance from the lens to the source, and \(D_S\) is the angular diameter distance from the observer to the source. These distances depend on both the source redshift \(z_S = 1.235\) and the unknown lens redshift \(z_L\).

By combining this equation with the Faber–Jackson relation (Faber & Jackson 1976) for elliptical galaxies, or the Tully–Fisher relation (Tully & Fisher 1977) for spiral galaxies, we generated a predicted observed magnitude for each filter as a function of lens redshift. The Faber–Jackson relation is given by

\[
M_B^0 = -19.4 + 5 \log h - 10(\log \sigma - 2.3)
\]

(de Vaucouleurs & Olson 1982), which becomes

\[
M_B = -18.9 + 5 \log h - 10(\log \sigma - 2.3),
\]

after applying the extinction correction \(B_T - B_V = 0.22\) (de Vaucouleurs et al. 1976), and with \(B_I = B_T + 0.29\) (Peterson et al. 1986). The Tully–Fisher relation is the same, but with \(\sigma\) replaced by circular velocity, which for an isothermal sphere is just \(\sqrt{2}\sigma\).

The predicted magnitude in the \(i'\) band is given by

\[
m_{i'} = M_B + DM(z_L) + K_{B,I}(z_L),
\]

where \(DM(z_L)\) is the cosmological distance modulus, and \(K_{B,I}(z_L)\) is the generalized \(K\)-correction between the lensing galaxy’s rest-frame \(B\) band magnitude and the observed \(i'\) band magnitude (see, e.g., Hogg et al. 2002). To calculate the \(K\)-correction for an elliptical galaxy at each potential lens redshift, we used a spectral energy distribution (SED) generated by the Bruzual & Charlot (2003) spectral evolution code. Our model consisted of a solar metallicity, instantaneous starburst at a redshift of 3.0, followed by passive evolution. For a spiral galaxy, we used an empirical Scd galaxy spectrum from Coleman et al. (1980), redshifted appropriately.

The observed magnitudes of the lensing galaxy matches those predicted for an elliptical galaxy for a range of redshifts \(0.4 \lesssim z \lesssim 1.0\). A spiral galaxy model also matches the \(H\) band observations, but would be brighter than observed at all redshifts in \(I\) by \(\gtrsim 1\) magnitude. The galaxy’s brightness and colors seem to be more consistent with an elliptical galaxy than a spiral.

In addition to this method, we estimated the probability distribution of the lensing galaxy’s redshift by calculating the lensing optical depth as a function of redshift, following the approach of Kochanek (1992). We found that the median redshift was 0.66, with a 68% confidence interval of \([0.41, 0.88]\). This is consistent with the results of the Faber–Jackson method.

### 6. CONCLUSIONS

The \(z = 1.235\) quasar HE 1113-0641 is lensed into a cross configuration, with four components ranging from 18.0 to 18.8 magnitudes in \(i'\). The maximum image separation is 0''67. A combination of ground-based and \textit{HST} imaging has
yielded reliable astrometry and photometry of the four quasar components of HE 1113-0641, as well as a good estimate of the position of the lensing galaxy. However, we were unable to measure the size or morphology of the galaxy. By assuming a circular pseudo-Gaussian shape and fixing a width for the galaxy, we were able to estimate its flux in the redder bands. A singular isothermal sphere (SIS) model succeeded in matching the positions of the quasar components and of the lensing galaxy, but was unable to match the observed flux ratios. Based on what has been seen with other gravitational lenses, it seems likely that this is due to perturbations from stellar microlensing or dark matter substructure in the lens galaxy.

The redshift of the lens galaxy remains elusive, but we estimate that \( z_L \sim 0.7 \pm 0.3 \), based on its observed flux and colors. At this redshift, its velocity dispersion, as measured by its lensing potential, would be 180 km s\(^{-1}\).

Because of its small separation, HE 1113-0641 will likely prove difficult to monitor using ground-based telescopes, and hope for measuring a spectroscopic redshift of the lensing galaxy is faint. Nevertheless, it is an interesting example of a small-separation lens, and may prove useful for studies that can take advantage of telescopes with very good seeing.

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