PATHSPACE DECOMPOSITIONS
FOR THE VIRASORO ALGEBRA
AND ITS VERMA MODULES

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Abstract
Starting from a detailed analysis of the structure of pathspaces of the \( \mathcal{A} \)-fusion graphs and the corresponding irreducible Virasoro algebra quotients \( V(c, h) \) for the \((2, q \text{ odd})\) models, we introduce the notion of an admissible pathspace representation. The pathspaces \( \mathcal{P}_\mathcal{A} \) over the \( \mathcal{A} \)-Graphs are isomorphic to the pathspaces over Coxeter \( \mathcal{A} \)-graphs that appear in FB models. We give explicit construction algorithms for admissible representations. From the finitedimensional results of these algorithms we derive a decomposition of \( V(c, h) \) into its positive and negative definite subspaces w.r.t. the Shapovalov form and the corresponding signature characters. Finally, we treat the Virasoro operation on the lattice induced by admissible representations adopting a particle point of view. We use this analysis to decompose the Virasoro algebra generators themselves. This decomposition also takes the nonunitarity of the \((2, q)\) models into account.

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1 Introduction

Ever since the paper [1], the close connection between CFT and statistical mechanics has been known. Subsequently, Belavin, Polyakov and Zamolodchikov [2] gave an in-depth analysis of the implications of conformal invariance in two dimensions [3, 4, 5]. The above mentioned correspondence can effectively be analyzed on the level of graphs and their pathspaces. These made their first appearance in statistical mechanics as the configuration spaces of integrable models defined by Andrews, Forrester and Baxter [6]. Since the subsequent identification of critical parameters by Huse [7] as belonging to the unitary discrete FQS-series of minimal models [8], there has been a desire to construct a Virasoro representation on these path spaces in order to better understand the appearance of this algebra in this context. This would be especially intriguing, since the algebra appears even in off-critical cases in the calculation of local height probabilities [9]. One interesting aspect would be a correspondence between the Temperly Lieb Jones algebra predominant in these models [10] and the Virasoro algebra itself providing a link between statistical systems and the corresponding conformal field theories [11, 12]. There has been a lot of research in developing new models along the lines of [6] and in the study of their properties [11, 13, 14, 15, 16, 17, 18, 19]. In [20] Forrester and Baxter introduced new models whose critical behavior is nonunitary [9]. In this context one can find the (2,q odd) series of minimal models in the critical behaviour of statistical models based on Coxeter $A_n$ graphs which will be studied in this paper. The class of models based on Coxeter graphs can be extended to include all A-D-E graphs [21] parallel to the classification of modular invariant partition functions [22, 23, 24, 25].

There are two ways in which to proceed in order to arrive at the correspondence mentioned above. One can start from the statistical model side or from the CFT side. There has been some development on the statistical side ([26] and ref. therein). In [26], the authors propose a double limit in order to obtain Virasoro generators from the Temperly Lieb Jones algebra. For a different approach see [27, 28, 29].

We shall tackle the above problem (construction of a Virasoro representation) from the CFT side where only little is known. On the CFT side, pathspaces appear in sum formulas for the conformal characters [30]. These pathspaces, however, are not over simple Coxeter graphs, but are more involved in the sense that the structure of the graphs which are considered is a bit more complicated. However, the $A_n$ graphs which appear in the (2,q) models (q odd) can be linked via pathspace isomorphisms to the Coxeter $A_n$ graphs [30] (see also section 2). In [31], another way of introducing graphs to CFT was found by rewriting character formulas.

In this paper, we will proceed from the pathspaces over the $A_n$ graphs. In section 2 we will give the relavant definitions and notations which we will need in the following. Then, in section 3, we will start the analysis of the pathspace and the Verma module quotient structure in order to establish a connection between them. As a result, we will introduce the notion of an admissible pathspace representation. In section 4, we will give and discuss possible constructions for these representations. The calculations involving these constructions will then lead us to the main conjecture which states signature character formulas for the $(2,q)$-models as well as the corresponding decomposition of the irreducible Virasoro quotient $V(c, h)$ into positive and negative definite subspaces $V(c, h) = V(c, h)^+ \oplus V(c, h)^-$. Finally, in section 5, we turn to the structure of the Virasoro action induced by admissible representations on the path space. To this end, a particle interpretation is given. The action of the Virasoro
generators $L_n$ can then be described by a shift ($\hat{L}_n$) and a one particle creation (resp. annihilation) ($C_n$) operator. The main conjecture can be restated in formulas for these operators and their adjoints. This reformulation explains the degree of nonunitarity of the $(2,q)$ or, in other words, the degree of non-selfadjointness of the Virasoro algebra in the pathspace metric. We conclude the paper with a summary of the results and an outlook into further fields of study.

2 Preliminaries and Notations

Since we are interested in representations of the Virasoro algebra, it is useful to first fix notations. By the Virasoro algebra $Vir$ we mean the complex Lie algebra generated by $L_n, n \in \mathbb{Z}$ and $C$ with commutation relations

$$ [L_n, L_m] = (m-n)L_{m+n} + \delta_{m+n,0}\frac{1}{12}C(m^2 - m) $$

$$ [L_n, C] = 0. $$

Please note the choice of signs in (2.1). As usual, we denote by $M(c, h)$ the Verma module to the eigenvalue $c$ of $C$ and lowest (choice of signs) weight $h$ for $L_0$. Furthermore, let $J(c, h)$ be the unique maximal submodule and $V(c, h)$ the unique irreducible quotient:

$$ V(c, h) = M(c, h)/J(c, h). $$

In this paper, we will concentrate on certain values of $c$, namely the ones belonging to the $(2,q)$ series. Here $(p, q)$ denotes the $c$ value:

$$ c(p, q) = 1 - \frac{6(p-q)^2}{pq}. $$

For these specific models, an explicit basis for the quotients $V(c, h)$ was given by Feigin, Nakanishi and Ooguri (FNO) in [32], the lowest weights being:

$$ h_j = -\frac{j(q-2-j)}{2q}, \quad j = 0, \ldots, N $$

with $q = 2N + 3$. The basis is given by elements

$$ L_{n_1} \cdots L_{n_m} |h_j\rangle $$

with $n_1 \geq \ldots \geq n_m \geq 1$ which satisfy the following two conditions:

1) $n_i - n_{i+N} \geq 2$ (difference two condition)

2) $\#\{n_i = 1\} \leq j$ (initial condition).

In the paper [32], the authors also introduce an order which will be quite helpful later on. One defines:

$$ L_{m_1} \cdots L_{m_r} \succ L_{n_1} \cdots L_{n_s} $$

if
i) \( r > s \),

ii) \( r = s \) and \( \sum_i m_i > \sum_i n_i \) or

iii) \( r = s, \sum_i m_i = \sum_i n_i \) and \( m_t > n_t, m_i = n_i \) for \( 1 \leq i \leq t - 1 \).

For the above basis a graphical enumeration method was given in [30]. The graphs used for this are the fusion graphs \( A_{N+1} \) of the corresponding theory with \( c(2N + 3, 2) \) (for a precise definition see [30]). Here we just define the graph \( A_N \) by its incidence matrix. For \( 0 \leq i, j \leq N - 1 \):

\[
(A_N)_{i,j} = \begin{cases} 
1 & \text{if } i + j < N \\
0 & \text{otherwise}
\end{cases}
\]

The basis itself is then given by paths over the corresponding graphs. By this we mean the following:

A path over a graph \( G \) with vertices numbered by a set \( I \) is a map \( : N_0 \mapsto I \) i.e. a sequence of vertices \( (l_i)_i \), with the restriction that two successive vertices are linked.

Furthermore, let \( P_{\{G, l_0, l_\infty\}} \) denote the Hilbert space with basis given by the paths on the graph \( G \) which start at \( l_0 \) and end in \( l_\infty \). This means that for \( i \gg 1, l_i = l_\infty \). The scalar product on the space is chosen to be the one in which all path are mutually orthonormal.

We then have a bijection [30] between \( P_{\{A_{N+1}, N-j, 0\}} \) and \( V(c(2, 2N+3), h_j) \) which is given by the simple map:

\[
((l_i)_i) \mapsto \ldots L_2^i L_1^i = \prod_{i=\infty}^{1} L_i^i \quad \text{for } i \in \mathbb{N}.
\] (2.8)

The restrictions for the paths given by the graph are directly translated into the difference two condition of the FNO basis. The index \( N + 1 \) fixes \( c \) to \( c(2, 2N+3) \) and the first term of the sequence \( l_0 \) dictates the \( h \) value of the theory and guarantees compliance with the initial condition. In the pathspaces we have an \( L_0 \) which provides the usual grading:

\[
L_0(l_i)_i := (h_{N-l_0} + \sum_{i \in \mathbb{N}} il_i) \cdot ((l_i)_i).
\] (2.9)

By K-theoretic arguments one can show that the pathspace over the \( A_{N+1} \) graph is isomorphic to the Coxter \( A_{2(N+1)} = A_{q-1} \) graph considered in [3] and [20].

**Example:** In the case of the Lee-Yang edge singularity (the (2,5) model), an isomorphism can easily be given. We just relabel the \( A_4 \) graph:

![Figure 1: Relabeling of the A4 graph](image)

1 2 3 4 1 0 0 1
and see that in this labeling the path space over $A_4$ is just the same as over the $A_2$ graph.

![Diagram](image)

Figure 2: $A_2$ fusion graph

One can also check that the $L_0$ (2.9) gives the same spectrum over the ground state as in the FB models. This example shows the obvious advantages of the $A$ graphs. The ground state $(\ldots 32323 \ldots)$ of the FB model reads simply as $(000 \ldots)$ and the operator $L_0$ just depends on the site $i$, not on three neighboring sites. We also have that the energy of the ground state is 0 and thus it is much easier to find the energy of the excited state, since no ground state contribution has to be subtracted.

In general, we can see that in the $A_{2n}$ models there are just $n$ distinct labels which can appear in the even resp. odd lattices, thus a graph with $N$ vertices should suffice. Furthermore, one feature of the $A$ graphs is that the vertex 0 is always connected to itself. So we always have a ground state $(000 \ldots)$ which replaces a ground state $(\ldots l, l - 1, l, l - 1 \ldots)$ and thus we also have a simpler structure for the excitations. In fact, part of this has already been realized in [6]. The authors introduce the following relabeling of the Coxeter $A_{r-1}$ graph, if $r$ is odd ($r = 2N + 3$):

$$l_i \mapsto \begin{cases} \frac{1}{2}(N + 1 - l_i) & \text{if } l_i \text{ odd} \\ \frac{1}{2}(l_i - 2) & \text{if } l_i \text{ even} \end{cases} \quad (2.10)$$

These labels correspond to the ones of the fusion graphs. Furthermore they map the initial conditions of the fusion and FB pathspaces for the $h$-values onto each other. Thus one can reassociate the fields of the given CFT to the edges of the graph in a procedure reverse to the one presented in [30].

In particular, for the $(2, 5)$ model we see that the fusion $A_2$ and the relabeled Coxeter $A_4$ graph are identical. In this way, all results for the pathspaces of the former can be translated back to the pathspace of the Coxeter $A_4$-graph. For the translation one must, however, first make a choice for the odd and even sublattices (see [6]).

The above relabeling was interpreted in [6] as a lattice gas picture. For the ABF resp. FB models, the restriction given by the graph amounts to the restriction that the sum of particles on neighboring lattice sites must be $N$ or $N - 1$. In this line of thinking, the fusion graphs describe a lattice gas with the perhaps more natural restriction that there are at most $N$ particles on two neighboring sites. The described particle interpretation as a lattice gas is adopted in this paper.

In the following we want to extend the operator $L_0$ (2.9) to a full representation of $Vir$.

## 3 Admissible Representations for $\mathcal{P}_A$
3.1 General Situation

We are looking for an irreducible representation of the Virasoro algebra on the pathspace $\mathcal{P}_A$. One other way to state this problem is given by the following:

**Remark:** If such a representation exists, then

$$\mathcal{P}_A \cong V(c, h) \quad \text{as Vir modules.} \quad (3.1)$$

Because of the universal property of the Verma module there exists a surjective homomorphism $\pi : M(c, h) \to \mathcal{P}_A$. Then $K := \ker(\pi)$ is a submodule. Therefore $K$ is a submodule of the unique maximal submodule $J(c, h)$.

Hence $K = J(c, h)$, since $\dim_n K := \dim_n J(c, h)$, using the isomorphism (2.8).

So the problem can be reformulated in the following way: The object we are looking for is an isomorphism

$$\Phi : \mathcal{P}_A \longrightarrow V(c, h). \quad (3.2)$$

One such isomorphism can be easily given in the following way:

**Example 1:**

$$\Phi((l_i)_i) := \prod_{i=\infty} L_i^i. \quad (3.3)$$

This is possible since the $\mathcal{A}$-graphs encode exactly the restrictions which appear for the FNO basis.

Example 1, however, does not respect all of the various natural structures of $\mathcal{P}_A$ and $V(c, h)$ which will be introduced in the following section.

3.2 Natural structures on $\mathcal{P}_A$ and $V(c, h)$

If we regard $\mathcal{P}_A$ (in the following denoted by $\mathcal{P}$) simply as a configuration space of pathspace origin, we can associate the following structures: First we have two operators $L_0$ and $K$ on $\mathcal{P}$, where $L_0$ is the usual pathspace $L_0$:

$$L_0((l_i)_i) := (h_{N-l_0} + \sum_{i \in \mathbb{N}} il_i) \cdot ((l_i)_i), \quad (3.4)$$

and $K$ is defined by:

$$K((l_i)_i) := (\sum_{i \in \mathbb{N}} l_i) \cdot ((l_i)_i). \quad (3.5)$$

Together they provide a double grading:

$$\mathcal{P} = \bigoplus_{n,k} \mathcal{P}_{n,k} \quad (3.6)$$

where

$$L_0|_{\mathcal{P}_{n,k}} = (n + h_{N-l_0}) \cdot id$$

and

$$K|_{\mathcal{P}_{n,k}} = k \cdot id.$$
Furthermore, we have a natural scalar product $( , )_P$ of the Hilbert space which takes the configuration space nature into account. It is defined by the fact that all paths are mutually orthonormal:

$$(((l_i)_i), ((g_i)_i))_P := \prod_i \delta_{l_i, g_i}.$$  (3.7)

This metric turns (3.6) into an orthogonal decomposition

$$\mathcal{P} = \perp_{n,k} \mathcal{P}_{n,k}.$$  (3.8)

So the pathspace basis should yield an orthogonal basis for the Verma module quotient under a reasonable representation (resp. isomorphism $\Phi$). And we should find a structure corresponding to the double grading (3.6).

If, on the other hand, we look at $V(c, h)$ we have the usual grading given by $L_0$ and a canonical hermitian form, the Shapovalov form (see e.g. [33]) which we denote by $\langle , \rangle_S$.

Together, these provide the usual orthogonal decomposition of $V(c, h)$:

$$V(c, h) = \perp_n V_n(c, h).$$  (3.9)

Although there is no obvious second grading which would provide another orthogonal splitting, we do have a filtration.

Consider the spaces

$$\tilde{V}_{n,k} := \text{span}\langle L_{n_1} \cdots L_{n_l} \mid n_i \geq \cdots \geq n_1, l \leq k \text{ and } \sum_i n_i = n \rangle$$

(i.e. linear combinations of products of at most $k \, L_i$’s of energy $n$.)

These spaces give the above mentioned filtration of each $\tilde{V}_n(c, h)$

$$\emptyset = \tilde{V}_{n,-1} \subset \tilde{V}_{n,0} \subset \cdots \subset \tilde{V}_{n,k-1} \subset \tilde{V}_{n,k} \subset \tilde{V}_{n,k+1} \subset \cdots \subset \tilde{V}_{n,\infty} = V_n(c, h).$$  (3.10)

The successive quotients of this filtration are especially interesting, since they have the same dimensions as the respective pathspace counterparts. To this end let $\hat{\mathcal{V}}_{n,k} := \tilde{V}_{n,k}/\tilde{V}_{n,k-1}$ denote these quotients then we have again using (2.8):

$$\text{dim} \hat{\mathcal{V}}_{n,k} = \text{dim} \mathcal{P}_{n,k}.$$  (3.11)

We denote by $\pi_{n,k}$ be the canonical projection $\pi_{n,k} : \tilde{V}_{n,k} \longrightarrow \tilde{V}_{n,k}$.

After these considerations it is clear which special features an isomorphism (3.2) should have in order to preserve the natural structures of $\mathcal{P}$. The following definition is therefore central.

We call an isomorphism $\Phi : \mathcal{P}_A \longrightarrow V(c, h)$ admissible, if the following two conditions hold:

a) $\Phi(\mathcal{P}_{n,k}) \subset \tilde{V}_{n,k}$ and thus $\pi_{n,k} \circ \Phi(\mathcal{P}_{n,k}) = V_{n,k}$.

b) $\Phi^*\langle , \rangle_S$ is orthogonal in the path-basis

i.e. $\langle \Phi((l_i)_i), \Phi((g_i)_i) \rangle = \epsilon((l_i)_i, (g_i)_i) \cdot ((l_i)_i, (g_i)_i)_P$ , with $\epsilon((l_i)_i, (g_i)_i) \in \mathbb{C}$. 

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Remarks:

1) We can always normalize so that the $\epsilon$ of b) is $\pm 1$.

2) The remark in condition a) encodes the information (3.11) about the dimensions of the corresponding spaces. Condition a) itself coarsely identifies the number of $L$'s with the sum of the nonzero entries in the path sequence. In particular, an entry $k$ signifies the presence of $k$ $L$'s (except for $l_0$, of course, which just fixes $h$).

3) Condition b) guarantees the existence of an orthogonal basis complementing the previous one in each step of the filtration. Reformulated, condition b) is equivalent to the existence of orthogonal splittings $s_{n,k}$ (i.e. $\pi_{n,k} \circ s_{n,k} = id$).

\[ 0 \rightarrow \tilde{V}_{n,k-1} \rightarrow \tilde{V}_{n,k} \xrightarrow{s_{n,k}} \hat{V}_{n,k} \rightarrow 0 \]  

(3.12)

which are orthogonal in the sense that

\[ s_{n,k}(\tilde{V}_{n,k}) \perp \tilde{V}_{n,k-1}. \]  

(3.13)

Together, these splittings give rise to an orthogonal decomposition: Let $V_{n,k} := s_{n,k}(\hat{V}_{n,k})$ then

\[ V(c,h) = \perp_{n,k} V_{n,k}. \]  

(3.14)

corresponding to the double grading (3.6) of the pathspace. The existence of such a basis, however, is nontrivial; it is rather the crucial point. If we look at example 1 (3.3) for instance, we see that it satisfies conditions a) while it fails to satisfy condition b).

In other words, the configuration space nature of the pathspace leads to the prediction of a specific orthogonal basis of $V(c,h)$.

4) An isomorphism $\Phi^\prime$ can be reconstructed from the splittings $s_{n,k}$ by:

\[ \Phi^\prime((l_i)_i) = s_{n,k} \circ \pi_{n,k}(\prod_{i=\infty}^1 L_i^h), \text{ for } (l_i)_i \in \mathcal{P}_{n,k} \]  

(3.15)

Although this isomorphism is not necessarily admissible, it satisfies the following two conditions

a) $\Phi^\prime(\mathcal{P}_{n,k}) \subset \tilde{V}_{n,k}$ and

b') $\Phi^\prime(\mathcal{P}_{n,k}) \perp \Phi^\prime(\mathcal{P}_{n',k'})$ w.r.t. $\langle \cdot, \cdot \rangle_S$, if $n \neq n'$ or $k \neq k'$.

An isomorphism satisfying the conditions a) and b') will be called weakly admissible. Given a weakly admissible isomorphism $\Phi^\prime$, we can always construct an admissible isomorphism $\Phi$ by orthogonalizing inside the spaces $\Phi^\prime(\mathcal{P}_{n,k})$. This is always possible, since the Shapovalov form is non-degenerate and by condition b') we already have an orthogonal decomposition $V(c,h) = \perp_{n,k} \Phi(\mathcal{P}_{n,k})$, so that the form is non-degenerate.
on the different parts of this decomposition; hence diagonalisable. Independent of the chosen orthogonalisation procedure we have:

\[ \Phi'(P_{n,k}) = \Phi(P_{n,k}) := V_{n,k}. \]  

(3.16)

So the spaces \( V_{n,k} \) as well as the decomposition corresponding to the double grading of the pathspace \( V(c, h) = \perp_{n,k} V_{n,k} \) depend only on the weakly admissible structure.

**Additional remark:** In some cases (if \( h = 0 \)), it is useful to restrict oneself only to the quasi-primary objects.

Let \( V_{n}^{q.p.} := \ker(L_{-1}|_{V(c,h)_{n}}) \),

\( \tilde{V}_{n,k}^{q.p.} := \ker(L_{-1}|_{\tilde{V}_{n,k}}) \),

resp. \( \hat{V}_{n,k}^{q.p.} := \ker(L_{-1}|_{\hat{V}_{n,k}}) \).

Analogously we define orthogonal splittings \( s_{n,k}^{q.p.} \) by:

\[ 0 \longrightarrow \tilde{V}_{n,k}^{q.p.} \longrightarrow \tilde{V}_{n,k}^{q.p.} \xrightarrow{s_{n,k}^{q.p.}} \hat{V}_{n,k}^{q.p.} \longrightarrow 0. \]  

(3.17)

The whole information can be retrieved due to the fact that the \( \text{SL}(2,\mathbb{C}) \) sub-Verma modules based on the \( V_{n}^{q.p.} \) provide a basis for \( V(c,h) \)

\[ \bigoplus_{n} U(L_{1})V_{n}^{q.p.} = V(c,h). \]  

(3.18)

The contravariance of the Shapovalov form and the fact that \( L_{1}|_{P_{n,k}} \in P_{n,k} \) if \( h = 0 \) (see section 4.1) guarantee that the orthogonal decomposition of the quasi-primary part

\[ V(c,h)^{q.p.} = \perp_{n,k} \hat{V}_{n,k} \]  

(3.19)

\( (\hat{V}_{n,k} := s_{n,k}^{q.p.}(V_{n,k}^{q.p.})) \) will induce the decomposition (3.14) on the whole of \( V(c,h) \).

This fact will be used for a construction procedure in the next section.

### 4 Admissible representations and signature characters

#### 4.1 Constructions for admissible representations

One way to find an orthogonal basis of \( V(c,h) \) (resp. the splittings \( s_{n,k} \)) with the required restrictions is the following:

**Construction 1:** We simply orthogonalize the basis vectors of FNO with respect to the FNO order (2.7).

As mentioned before, we do not know a priori that this algorithm will work, since the Shapovalov form is not definite and isotropic vectors may occur. But on the other hand, the
existence (resp. the success of the algorithm) will be exactly the new information gained. The corresponding isomorphism would be:

$$\Phi_{\text{FNO}}(l_i) := \left( \prod_{i=1}^{\infty} L_i^{l_i} \right)^{\perp_{\text{FNO}}}$$

where the superscript $\perp_{\text{FNO}}$ refers to the above mentioned orthogonalization. In doing the calculations, we proceeded as follows: We just generated the Shapovalov form in the FNO basis and used it to do the orthogonalization. This was carried out up to $V_{32,5}$ (that is for a 93 dimensional subspace) for the $(2,5)$ model and up to $V_{15,5}$ (that is for a 37 dimensional subspace) for the $(2,q$ odd) models for $q < 35$.

We could also use any other order in which:

$$L_{m_1} \cdots L_{m_r} \succ L_{n_1} \cdots L_{n_s}$$

if $r > s$. The FNO order is, however, best suited (see section 4.2).

**Construction 2:** As mentioned before, if $h = 0$, we can restrict ourselves to the quasi-primary objects. Hence, we first enlarge a basis of $\tilde{V}^{q,p}_{n,k-1}$ which we already know by induction to a basis of $\tilde{V}^{q,p}_{n,k}$ and orthogonalize the new vectors in an arbitrary fashion. This procedure can be refined in the sense that the FNO ordering provides an even finer filtration as (4.10) (as above, we could choose any other ordering with the condition (4.2)), so we can order the new basis and orthogonalize w.r.t. this ordering as in construction 1. Although this construction seems more tedious, especially if one wants to recover the full basis, it can be quite useful for the calculation of the signature, since all $L_1$ descendent vectors of quasi-primary ones have the same sign of the metric. The calculation itself was performed in the following manner: If $h=0$ we have the nice feature that each singular vector can be taken to be of the form:

$$L_{m_1}L_{m_2}L_{n_r} \cdots L_{n_1} + \text{lower order terms w.r.t. the ordering (2.7).}$$

If $k = 2$ this is especially nice since $V^{q,p}_{n,2}$ is onedimensional. The construction was carried out for $k = 2$ up to $n = 200$ for all $(2,q)$ models.

We have used both these constructions explicitly up to different grades for $n,k$. The results of the algorithms and the terms of the resulting signature characters are contained in the next section.

### 4.2 Signature characters for the $(2,q)$ models

Since any admissible representation results in an orthogonal decomposition and, in addition, provides an orthogonal basis of $V(c,h)$ (corresponding to the path basis), it can be used to calculate signatures of the different parts of the decomposition (3.14). As mentioned above, this has been carried out for various values of $n$ and $k$.

For all calculated examples we find that:

$$\text{sign}(\langle , \rangle)_{V_{n,k}} = (-1)^k \text{dim} V_{n,k}.$$  

In other words, the Shapovalov form is $(-1)^k$ definite on $V_{n,k}$. This is truly a remarkable result: from the pathspace nature we do not only get an orthogonal basis, but we also find,
and this is very important, a basis which provides the splitting of $V(c, h)$ into positive and negative definite subspaces:

$$V(c, h) = V(c, h)^+ \oplus V(c, h)^-.$$  \hfill (4.5)

For a discussion of the alternating sign of the definiteness see remark in section 5.3.

The data presented are far out of the region of mere coincidence and thus lead us to the following conjecture:

**Main conjecture:** In the $(2, q)$-minimal models admissible representations exist, and they provide a splitting (4.5) of $V(c, h)$ into positive and negative definite subspaces by

$$V(c, h)^+ := \bigoplus_{k \text{ even}} V_{n,k}, \quad V(c, h)^- := \bigoplus_{k \text{ odd}} V_{n,k}.$$  \hfill (4.6)

The resulting signature character would be:

$$\sigma(c(2, 2N + 3), h_j)(q) = \sum_{n_1, \ldots, n_N \geq 0} (-1)^{\sum_{i=1}^{N_i} in_i} \frac{q^{N_1^2 + \cdots + N_k^2 + \cdots + N_N} (q)_{n_1} \cdots (q)_{n_N}}{(q)_{n_1} \cdots (q)_{n_N}}$$  \hfill (4.7)

with $N_i = \sum_{j=1}^{N_i} n_j$ and $(q)_i = \prod_{j=1}^{i} (1 - q^j)$.

**Remarks:**

1) The signature character formulas (4.7) have been previously conjectured by Nahm [34]. They have been compared to the ones given by Kent [35, 36] up to $O(q^{100})$ for the $(2, 5)$ and the $(2, 7)$ model and the two coincide [37].

2) The exact structure of the splittings or of the orthogonalization is of no importance for the induced metric on $P$, since it is positive resp. negative definite on each of the $P_{n,k}$. So any other choice of $\Phi$ would yield the same results. Only the weak admissibility is of importance.

3) The characters corresponding to (4.7) are the same except for the term $(-1)^{\sum_{i=1}^{N_i} in_i}$ and have been given in [32]. They correspond to a series of sum rules for characters which can be interpreted in terms of quasiparticles [31]. In this setting, the exponent $\sum_{i=1}^{N_i} in_i$ is just the sum over the number of quasi particles of type $i$ weighted with their energy weights (see also section 5.1).

## 5 Virasoro on the lattice

### 5.1 General remarks

Any isomorphism $\Phi$ (3.2) is, of course, nothing but a realization of the Virasoro algebra on the lattice in the sense of section 1. If we now look at admissible representations only, we can describe the action on the paths themselves more explicitly. The total structure is, however, dependent on the specific choice of basis induced by the choice for $\Phi$. In general,
when considering only the weakly admissible structure of an admissible representation, we can find the following:

\[ L_m(\tilde{V}_{n,k}) \subset \tilde{V}_{n+m,k+1} \quad \text{and} \quad (5.1) \]
\[ L_{-m}(\tilde{V}_{n,k}) \subset \tilde{V}_{n-m,k}. \quad (5.2) \]

Furthermore, we see from the restrictions for the annihilating ideals \[32\]:

\[ L_1^l(\tilde{V}_{n,k}) \subset \tilde{V}_{n+l,k+j}, \quad (5.3) \]
\[ L_2^r(\tilde{V}_{n,k}) \subset \tilde{V}_{n+2r,k+N+j}, \quad (5.4) \]

for \( c = c(2, 2N + 3) \) and \( h = h_j \).

The first inclusion is given by the initial condition and the second one by the difference two condition.

The inclusions (5.1) and (5.2) are most easily seen in the FNO basis:

\[ L_m(L_{n_1} \cdots L_{n_k}) = \sum_{i=k}^{r} [(n_i - m) L_{n_k} \cdots L_{n_{i+m}} \cdots L_{n_1}] \]
\[ + L_{n_k} \cdots L_{n_{r+1}} L_m L_{n_r} \cdots L_{n_1}, \text{where } n_{r+1} > m \geq n_r. \]

If any of the above vectors violates the conditions for the FNO basis, its expression in terms of the basis vectors contains not more \( L_i \)'s than before. This is seen directly from the annihilating ideals and the ordering as in the proof in \[32\].

If we now pull the Virasoro action back with an admissible \( \Phi \) onto the pathspace, the above inclusions translate directly into restrictions on the \( P_{n,k} \)

\[ L_m(P_{n,k}) \subset P_{n+m,k} \oplus P_{n+m,k+1} \quad (5.6) \]
\[ L_{-m}(P_{n,k}) \subset P_{n-m,k-1} \oplus P_{n-m,k}. \quad (5.7) \]

This can be seen in the pulled back Shapovalov form: First of all, we have the inclusions (5.1) and (5.2) for the path spaces as well

\[ L_m(P_{n,k}) \subset \bigoplus_{j=1}^{k+1} P_{n+m,j} \quad \text{and} \quad (5.8) \]
\[ L_{-m}(P_{n,k}) \subset \bigoplus_{j=1}^{k} P_{n-m,j}. \quad (5.9) \]

Furthermore we see from the orthogonality (3.14):

\[ \Phi^* \langle (l_i)_i, L_m((g_i)_i) \rangle_S = \langle \Phi((l_i)_i), L_m(\Phi((g_i)_i)) \rangle_S = \langle L_{-m}\Phi((l_i)_i), \Phi((g_i)_i) \rangle_S = 0 \]

by (3.13), since \( L_{-m}\Phi((l_i)_i) \in \tilde{V}_{n,k-1} \) and \( \Phi((g_i)_i) \in s_{n,k}(\tilde{V}_{n,k}). \)
So (5.6) follows. (5.7) follows in the same manner. From the initial condition (5.3) and from the difference two condition (5.4) we find by similar arguments:

\[ L_l \subset P_{n,k} + l,k \oplus \ldots \oplus P_{n+l,k+j}, \quad \text{and} \quad (5.11) \]

\[ L_r \subset P_{n+2r,k} + 2r,k \oplus \ldots \oplus P_{n+2r,k+N+j}, \quad (5.12) \]

if \( c = c(2,2N+3) \) and \( h = h_j \).

These restrictions concerning the action on the paths lead us to the following additive splitting of the usual Virasoro generators:

Let \( L_m = \hat{L}_m + C_m, \text{ for } m \in \mathbb{Z} \) \hspace{1cm} (5.13)

where \( \hat{L}_m|_{P_{n,k}} = P_{n,k} \circ L_m \)

\( C_m|_{P_{n,k}} = P_{n,k+1} \circ L_m, \text{ for } m > 0 \)

\( C_m|_{P_{n,k}} = P_{n,k-1} \circ L_m, \text{ for } m < 0 , \)

where the \( P_{n,k} \) denote the orthogonal projection onto \( P_{n,k} \). The relations (5.3) and (5.4) can now be written as:

\[ C_{j+1}^1 = 0 \]

\[ C_{N+j+1}^2 = 0, \text{ if } c = c(2,2N+3) \text{ and } h = h_j. \] \hspace{1cm} (5.14)

(5.15)

This means that the information about \( h \) resp. \( c \) is encoded in the operators \( C_1 \) resp. \( C_2 \). These characteristic quantities can now be simply read off from the nilpotency index of the respective operator.

We can also find commutation relations for these operators by substituting (5.13) into the basic Virasoro relation (2.1) and then comparing the degrees of the various operators w.r.t. the \( K \) grading.

To simplify things, consider from now on \( \hat{L}_n \) and \( C_n \) for \( n \in \mathbb{N} \) only and \( \hat{L}_n^\dagger = \hat{L}_{-n} \) resp. \( C_n^\dagger = C_{-n} \), where the dagger is the adjoint w.r.t. the Shapovalov form.

We then have (as well as the resp. daggered equations):

\[ [\hat{L}_n, \hat{L}_m] = (m-n)\hat{L}_{n+m} \]

\[ [\hat{L}_n, C_m] + [C_n, \hat{L}_m] = (m-n)C_{n+m} \]

\[ [C_n, C_m] = 0 \]

and as mixed commutators if \( m \geq n \):

\[ [\hat{L}_n^\dagger, \hat{L}_m] + [C_n^\dagger, C_m] = (m+n)\hat{L}_{m-n} + \delta_{m,0} \frac{c}{12} (m^3 - m) \]

\[ [\hat{L}_n^\dagger, C_m] = (m+n)C_{m-n} \]

\[ [C_n^\dagger, \hat{L}_m] = 0 \]

and as mixed commutators if \( m \leq n \):

\[ [\hat{L}_n^\dagger, \hat{L}_m] + [C_n^\dagger, C_m] = (m+n)\hat{L}_{n-m} + \delta_{m,0} \frac{c}{12} (m^3 - m) \]

\[ [\hat{L}_n^\dagger, C_m] = (m+n)C_{n-m} \]

\[ [C_n^\dagger, \hat{L}_m] = 0 \]
Although these relations seem a bit more complicated than the original Virasoro relations, there are several reasons for the introduction of the operators $\hat{L}_n$ and $C_n$. One is that in a particle interpretation of the path space they can be viewed as shift resp. creation-annihilation operators. There are several ways in which to give a particle interpretation for the sequences $(l_i)$. In [31] for instance, a quasi particle interpretation was developed. According to this any $(2,q)$ model contains $N$ different types of quasi particles.

Here we shall adopt a simpler point of view. A nonzero $l_i$ is just taken to signify the presence of $l_i$ particles on the site $i$. The connection to the quasi particle picture is just looking at a quasi particle of type $i$ as being made up by $i$ single particles. In this context, the spaces $P_{n,k}$ are the configuration spaces of $k$ particles of total energy $n$. In these terms the $C_n$ are one particle creation operators and the $C_n^{\dagger}$ one particle annihilation operators, while the $\hat{L}_n, \hat{L}_n^{\dagger}$ conserve the total number of particles and thus just produce a net number of $n$ shifts to the right resp. to the left. In particular, we have

$$[\hat{L}_1, \hat{L}_n] = (n - 1)\hat{L}_{n+1} \quad \text{for} \quad i \in \mathbb{N} \quad (5.25)$$

$$[\hat{L}_1, C_n] + [C_1, \hat{L}_n] = (n - 1)C_{n+1}. \quad (5.26)$$

This expresses that higher order shifts are obtained by nested $\hat{L}_1$ commutators of $\hat{L}_2$. Furthermore, the commutators simplify for $h = 0$. Then we have from the initial condition that:

$$C_1 = C_{-1} = 0, \quad \text{thus} \quad \hat{L}_1 = L_1 \quad (5.27)$$

and the above relations simplify to

$$[L_{\pm 1}, \hat{L}_n] = (n \mp 1)\hat{L}_{n\pm 1} \quad (5.28)$$

$$[L_{\pm 1}, C_n] = (n \mp 1)C_{n\pm 1}. \quad (5.29)$$

Now higher order shifts and creation (resp. annihilation) of particles by Virasoro action are obtained by nested $L_1$ commutators of $\hat{L}_2$ and $\hat{C}_2$ (resp. $C_{-2}$). The above remarks characterize the net action of the Virasoro on a $k$-particle state. The finer structure -how the shifts look on the single particles or where along the chain a new particle is created- depends on the specific choice of basis for the isomorphism $\Phi$. In the next section we will discuss this question for the constructions 1 ((4.1) i.e. the orthogonalized FNO basis) and 2 (quasi-primary construction).

### 5.2 Shifts and Creation

The simplest action of $L_n$ one could imagine would be that $\hat{L}_n$ shifts each particle $n$ sites and $C_n$ creates a particle at site $n$. This is almost the case in example 1. An action like this would be something like the Sugawara construction [33]. Problems arise, however, if the new configuration is not allowed. Here the action becomes very complicated. It can happen, for instance, that all particles are shifted or more than one particle is annihilated etc. This is an effect induced by the nonadmissibility of this example. The admissibility guarantees an overall control over the action as discussed in the previous section. The price paid is the "loss" of a simple shift and creation structure. In construction 1, however, we have further control over the Virasoro action. To this end we associate to a sequence $(l_i)$ its finite number of nonzero members:

$$(l_i)_i \rightarrow (l_{i_1}^1, \ldots, l_{i_m}^m) \quad (5.30)$$
that is the configuration with $l_k$ particles at site $i_k$.
We now examine the specific action of the Virasoro for construction 1:

1) $\hat{L}_n$ results in a sum of configurations where the highest particle is shifted at most $n$ steps. In the ordering (2.7) we have:

$$\hat{L}_n(n_{i_1}, \ldots, n_{i_m}) = \sum_{\text{config.} \in \mathcal{P}_{n_k}} \lambda_{\text{config.} \text{config.}}$$

with $\sum_i n_i = k$ and $\text{config.} \geq (n_{i_1}, \ldots, (n_m - 1)^{i_m}, 1^{i_m+n})$.

2) $C_n$ creates a sum of configurations, in which the new particle is at most at site $n$:

$$C_n(n_{i_1}, \ldots, n_{i_m}) = \sum_{\text{config.} \in \mathcal{P}_{n_{k+1}}} \lambda_{\text{config.} \text{config.}}$$

with $\sum_i n_i = k$ and

$$\text{config.} \leq \begin{cases} (n_{i_1}, \ldots, n_{r+1}^i, 1^r, n_{r+1}^i, \ldots, n_{i_m}) & \text{if } r < n < r+1 \\ (n_{i_1}, \ldots, (n_r + 1)^i, \ldots, n_{i_m}) & \text{if } r = n. \end{cases}$$

If any of the configurations above is not allowed, the respective coefficient is zero. These relations are proven analogously to the inclusions (5.1) and (5.2). Of course, there are analogous restrictions for the daggered operators.

**Example:** In particular in the vacuum sector of the (2,5) model $C_2$ creates only particles at site 2. So, for instance,

$$C_2(1^2, n_{i_2}^2, \ldots, n_{i_m}^m) = 0$$
$$C_2(1^3, n_{i_2}^2, \ldots, n_{i_m}^m) = 0$$
$$C_2(1^4, n_{i_2}^2, \ldots, n_{i_m}^m) = \sum_{\text{config.} \in \mathcal{P}_{n_{k+1}}} \lambda_{\text{config.} \text{config.}}$$

with $\text{config.} \leq (1^2, 1^4, \ldots, n_{i_m}^m)$

resp.

$$C_2^\dagger(1^2, n_{i_2}^2, \ldots, n_{i_m}^m) = \sum_{\text{config.} \in \mathcal{P}_{n_{k+1}}} \lambda_{\text{config.} \text{config.}}$$

with $\text{config.} \geq (n_{i_2}, \ldots, n_{i_m}^m)$

$$C_2^\dagger(1^3, n_{i_2}^2, \ldots, n_{i_m}^m) = 0$$
$$C_2^\dagger(1^4, n_{i_2}^2, \ldots, n_{i_m}^m) = 0.$$

The relevant coefficients of this example have been calculated up to $n=16$ and comply with the above restrictions.

If we now turn to construction 2, we see that the action of $L_1$ is particularly simple. In fact one defines the whole isomorphism by the following action:

$$L_1(n_{i_1}^i, \ldots, n_{i_m}^m) = (n_{i_1}^i, \ldots, (n_m - 1)^{i_m}, 1^{i_m+1}).$$
The rest of the Virasoro action is, however, very complicated, since there are no further restrictions apart from (5.6) and (5.7), so that in the generic case all coefficients $\lambda_{\text{config}} \neq 0$. In fact, we have not found any other admissible representations with harder restrictions on the overall action of $\text{Vir}$ than construction 1. For constructions resulting from a different ordering this is fairly easily seen.

5.3 Nonunitarity and the pathspace metric

If we regard the adjoint $\dagger$ with respect to $(\ , \ )_P$ we have no nice formulas for $L_n$. This is a general fact which results from the negative $c$ values. It has been remarked several times [9, 28, 29, 33] that one can define a positive definite scalar product on the configuration space, but that in this product the Virasoro algebra is not self-adjoint. In fact, it cannot be due to the negative $c$. The complex structure of the shifts can also be explained in this way. For the treated $c$ values one cannot find a Sugawara construction in which the Heisenberg algebra and the Virasoro algebra are both self-adjoint [33]. From the main conjecture, we have however:

$$\hat{L}_n^\dagger = \hat{L}_{-n} \quad \text{ (5.34)}$$

$$C_n^\dagger = -C_{-n}. \quad \text{ (5.35)}$$

These equations establish the connection between the positive definite scalar product on the configuration space and the way the Virasoro algebra behaves under adjungation. In this way, the above relations can also be taken to be a reformulation of the main conjecture. In the particle language, this reads in the following way: The nonunitarity of the $(2,q)$ models manifests itself only in the different sign between the creation and annihilation operators under adjungation.

**Remark:** In a way this is the simplest possible nonunitarity. If we want to have the pulled back Shapovalov form to be definite (positive or negative) on the configuration space of $k$ particles $\oplus_k P_{n,k}$ then the definiteness changes sign from the space of $k$ particles to the space of $k+1$ particles. Let $(l_i)_i \in P_{n,k}$ and $m \gg n$. From our assumption of definiteness we have (5.34) and $C_n^\dagger = \pm C_{-n}$ so:

$$((l_i)_i, [L_{-m}, L_m](l_i)_i)_P = \frac{2mn + c}{12}(m^2 - m) < 0 \quad \text{ (5.36)}$$

$$= ((l_i)_i, L_{-m}L_m(l_i)_i) \quad \text{ (5.37)}$$

$$= ((l_i)_i, \hat{L}_{-m}\hat{L}_m(l_i)_i)_P + ((l_i)_i, C_{-m}C_m(l_i)_i)_P \quad \text{ (5.38)}$$

$$= (\hat{L}_m(l_i)_i, \hat{L}_m(l_i)_i)_P \pm (C_m(l_i)_i, C_m(l_i)_i)_P, \quad \text{ (5.39)}$$

so the second term must be negative, so we must have $C_n^\dagger = -C_{-n}$.

6 Conclusion and outlook

We investigated the action of the Virasoro algebra on the pathspace given by the $A_{N+1}$ graphs. These spaces are related by a pathspace isomorphism to the configuration space of the FB models on coxeter $A_{q-1}$ graphs whose critical indices are those of the $(2,q)$ models ($q = 2N+3$). A detailed study of the spaces under consideration leads us to the central notion
of an admissible representation. Explicit constructions for such representations were given. The finite dimensional results concerning these representations lead us to conjecture signature character formulas for the above stated models as well as the respective decomposition of $V(c, h)$ into positive and negative definite subspaces. Furthermore, the explicit Virasoro action corresponding to an admissible representation was studied. As a result, we find that the action of $L_n$ can be described by a shift ($\hat{L}_n$) and a creation ($C_n$) operator. The parameters $N$ and $j$ which specify the central charge $c$ and the lowest weight $h$ of the theory appear as nilpotency indices for these operators. Finally, it was shown that the degree of nonunitarity resp. of the non-selfadjointness of the Virasoro algebra of the $(2, q)$ models can be understood in terms of the formulas for the adjoint operator of $\hat{L}_n$ and $C_n$.

These results establish a close connection between integrable models and the CFTs considered here. On the level of the graphs the pathspace isomorphism between the fusion graph $A$ explains the appearance of the $(2, q)$-series in the $A_{q-1}$ models. From the interplay between the two perceptions of the objects as belonging to CFT or statistical mechanics we gain insight into their structure. From the CFT side we learn how the Virasoro algebra operates as a sum of shifts and creation. The integrable model side provides us with signature characters and their corresponding decompositions into positive and negative definite subspaces.

We hope that in this spirit we can learn more about the connection between CFT and integrable models. There are also other series of CFT’s in which graphs appear in sum rules for characters [31]. A similar treatment would be of interest. Perhaps there is also a connection to the quasi particle interpretation of sum rules [38, 39, 40]. In this way, one can maybe also understand the implications of the simple structure of the signature characters and their role in statistical mechanics.

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