Effect of oxygen pressure on the structural and optical properties of BaSnO₃ films prepared by pulsed laser deposition method.

JOHN, J., CHALANA, S.R., PRABHU, R. and PILLAI, V.P.M. 2019

This is a post-peer-review, pre-copyedit version of an article published in Applied Physics A. The final authenticated version is available online at: http://dx.doi.org/10.1007/s00339-019-2432-0.
Effect of Oxygen pressure on the structural and optical properties of BaSnO₃ films prepared by pulsed laser deposition method

Jibi John¹, Chalana S. R¹, Radhakrishna Prabhu² and V. P. Mahadevan Pillai¹*

¹Department of Optoelectronics, University of Kerala, Kariavattom 695581, Thiruvanthapuram, Kerala, India.
²School of Engineering, Robert Gordon University, Aberdeen, UK.

* E-mail: vmpillai9@gmail.com

Abstract

BaSnO₃ thin films were prepared on quartz substrate by pulsed laser deposition technique under different background oxygen pressure. The effects of oxygen pressure on the structural, morphological and optical properties of BaSnO₃ thin films were investigated using techniques such as X-ray diffraction (XRD), Scanning Electron Microscopic image analysis (SEM), energy-dispersive x-ray (EDX) spectroscopy, Ultraviolet–Visible (UV–Vis) spectroscopy and X-ray Photoelectron Spectroscopy (XPS). All the films were annealed at 700 °C for 1hr and the laser energy used is 60 mJ. XRD analysis reveals that there is a systematic increase of intensity of <110> with increase in oxygen pressure up to 0.02 mbar and thereafter shows reduction in intensity of the peak. A SEM and AFM image also reveals that the 0.02 mbar oxygen pressure has good morphology. XPS reveals that the presence of constituent elements in the prepared sample. The prepared sample has high transmittance in the visible region and infrared region.

Introduction

Perovskites with the general formula MSnO₃ (M = Ba, Sr, Ca and Pb) have found wide variety of applications [1]. Barium stannate (BaSnO₃) is a cubic perovskite oxide compound [2] that behaves as an n-type semiconductor [3] with a wide band gap of 3.4 eV [3-6]. BaSnO₃ has three-dimensional framework of corner-sharing SnO₆ octahedra, having a linear Sn-O-Sn bonding angle 180° with ideal cubic Pm3m space group [7]. By doping, with suitable dopants materials the properties of the BaSnO₃ can be modified in a wide range due to its high substitutional flexibility [8,9]. BaSnO₃ is used to prepare capacitors because of its characteristic dielectric properties [10]. It is also used as a sensor material for a wide variety of gases, such as CO, Cl₂, NOₓ, O₂, CO₂, C₂H₅OH, CH₃SH and sensing humidity and liquefied petroleum gas [11,12]. The preparation of BaSnO₃ powder by conventional solidstate reaction through sintering BaCO₃ and SnO₂ at a high temperature ranging from 1,200 °C to 1,400 °C was reported by S. Upadhyay
et al. [13,14]. Hydrothermal synthesis route for the preparation of BaSnO$_3$ is reported by Kutty and Vivekanadan [15]. Mildly lanthanum doped BaSnO$_3$ exhibits very high electrical mobility of $320 \text{ cm}^2(\text{Vs})^{-1}$ at room temperature and superior thermal stability at high temperatures [16]. Lampe et al. [11] reported the NO and CO sensing behaviour of BaSnO$_3$ thin films at the temperature range of 450–550 °C, and of 600–700 °C respectively [17]. The preparation of lanthanum doped BaSnO$_3$ with substrate temperature 850 °C and oxygen partial pressure $1 \times 10^{-3}$ Torr using pulsed laser deposition is reported by Sallis et al. [18]. A carrier concentration of $2 \times 10^{21}$ cm$^{-3}$ and mobility 0.69 cm$^2$/Vs was reported in lanthanum-doped barium stannate films [19]. Although there are different techniques reported for thin film preparation the pulsed laser deposition (PLD) method is simple and unique technique which can be used to prepare good quality films from a large amount of materials and compounds [20]. In this method one can vary substrate temperature, background pressure, substrate to target distance, laser fluence, etc.. PLD is one of the best methods for preparing the films which maintains the composition close to that of target [20]. There are a large amount of studies in bulk and doped BaSnO$_3$ powder, but studies in thin films of BaSnO$_3$ are very rarely reported. Thin films of iron and lanthanum doped barium stannate is reported by Jayaraj et al and its magnetic properties are studied [21,22]. The electrical properties of BaSnO$_3$ doped lanthanum thin films by radio frequency magnetron sputtered method are reported by Lou et al [23]. This paper reports BaSnO$_3$ thin films prepared by pulsed laser deposition technique and a study on the effect of background oxygen pressure on its structural, morphological and its optical properties.

2. Experiment

BaSnO$_3$ powder was prepared by the conventional high temperature solid state method [13,14]. High quality (purity-99.99%, Sigma) BaCO$_3$ and SnO$_2$ powders are mixed in stoichiometric ratio and the mixture was ground well in an agate mortar for 6 hours using acetone as the mixing agent. The grinded powder taken in a platinum crucible was heated in a microprocessor programmable furnace for 6 h at a temperature of 1250°C. During heating the rate of rise of temperature was kept at 5°C/min. After heating, the powder was cooled to room temperature naturally. The phase purity of the BaSnO$_3$ powder thus prepared was analyzed by the X ray diffraction pattern and micro-Raman spectrum. The BaSnO$_3$ powder thus prepared was well-ground in an agate mortar and poly vinyl alcohol was added to the mixture as the binding agent. Pellets of 3 mm diameter and 1 mm thickness were prepared from the BaSnO$_3$ powder in a hydraulic press by applying pressure of 8 ton and the pressed pellet was sintered at temperature of 1000°C for 6 h in a programmable furnace. The well-sintered pellet thus prepared is used as the target for film preparation. The
deposition of BaSnO$_3$ films was done in a multiport chamber which was pre-evacuated to a base pressure of $10^{-6}$ mbar using a turbo molecular pump and a rotary pump. The ablation of target was carried out using a laser radiation at 355 nm wavelength (pulse width 7 ns and repetition rate 10 Hz) from a Q switched (frequency tripled) Nd:YAG laser (Quanta Ray INDI Series-Spectra Physics). Cleaned quartz substrate kept at distance of 4.5 cm from the target was heated at a temperature of 600°C. The films are fabricated on heated substrate for a deposition time of 45 min using a laser fluence of 60 mJ at different background oxygen pressures viz 0, 0.01, 0.015, 0.02, 0.025 and 0.03 mbar. To maintain uniform ablation during film preparation the target was rotated (speed 33 rpm). All the films are annealed at temperature of 700°C for 1 h. The film prepared without oxygen ambience is designated as BS0 and the films prepared with background oxygen pressures 0.01, 0.015, 0.02, 0.025 and 0.03 mbar are designated as BS 0.01, BS0.015, BS0.02, BS0.025 and BS 0.03 respectively.

Crystalline structure of the BaSnO$_3$ films was characterized by XPERT PRO diffractometer in the 20 range 20-80° using Cu-K$_{\alpha_1}$ radiation of wavelength 1.5406 Å in Bragg Brentano geometry. The measurements were carried out with a step size of 0.04000 and a scan speed of 2° min$^{-1}$ was maintained. Surface morphology of the films was analyzed using Atomic Force Microscopy (AFM) and Scanning Electron Microscopy (SEM) at nanometric scale using by Bruker, Dimension Edge with ScanAssist in contact mode and Nova Nano SEM-450 Field Emission Scanning Electron Microscope (FEI-USA) respectively. The thickness of the films was measured using lateral SEM analysis. Photoluminescence spectra of the films were recorded using Horiba Jobin Yvon Fluorolog III modular spectrofluorometer equipped with Xe-flash lamp using an excitation wavelength of 350 nm. The optical absorption, transmission spectroscopic measurements of all the films were performed in the wavelength range of 2300–250 nm using Perkin Elmer Lambda 950 spectrophotometer. The surface composition and elemental analysis were done using X-ray Photoelectron Spectrometer (XPS) ESCALABX1+A1528 using avantage software.

Results and discussion

XRD analysis

XRD patterns of the laser ablated BaSnO$_3$ films are shown in Fig. 1 which indicate the polycrystalline nature of the films and all the XRD patterns present peaks corresponding to cubic BaSnO$_3$ crystalline phase [JCPDS Ref. Code 15-0780]. The XRD pattern of BS 0 film presents two high intense peaks at 20 values 30.76° and 44.04°, a medium
intense peak at 54.67 ° and weak peak at 2θ value 72.72°. These peaks can be indexed to lattice reflection planes (110), (200), (211) and (300) of cubic BaSnO₃ crystalline phase respectively. In addition to these four peaks, a medium intense peak at 2θ value 64.05° corresponding to (220) lattice reflection plane of cubic phase can be seen in the XRD patterns of the films prepared under background oxygen. In the XRD patterns of all the films, (110) plane possesses the highest intensity indicating <110> as the preferred direction of crystalline growth in films. In films prepared under background oxygen pressure (ρO₂), the intensity of (110) peak shows a systematic increase with increase in ρO₂ up to 0.02 mbar and thereafter the intensity of the peak decreases with increase in ρO₂ (Fig. 2a). The β(hkl), Full Width at Half Maximum (FWHM), of the (110) peak is estimated from the XRD data and are given in Table 1 and for the BSO film its value is 0.262 °. The value of β(hkl) of (110) peak shows a systematic decrease with increase in background oxygen pressure up to ρO₂ = 0.02 mbar and thereafter it increases with increase in background oxygen pressure. From this it can be inferred that moderate oxygen ambience will be helpful in getting good crystalline BaSnO₃ film and the film prepared at ρO₂ = 0.02 mbar possesses the highest crystalline quality.

From the XRD patterns of the films, it can be seen that as ρO₂ increases, both (110) and (200) peaks show large intensity enhancement up to ρO₂ = 0.02 mbar. The relative intensities ρ(110) and ρ(200) of the peaks (110) and (200) for the films prepared under background oxygen pressure with respect to those values in BSO films are estimated and are plotted in the Fig. 2b. This indicates that at low or moderate oxygen pressure there exists a competition for growing the crystallites along <110> and <200> directions. At higher oxygen pressures these competing tendency of crystal growth disappears and the crystallite prefer to grow along <110> direction. The observed deterioration in the crystalline of the films at higher background oxygen pressures (> 0.02 mbar) may be due to the decrease in the deposition rate with increase in the background oxygen ambience [24]. Beyond 0.02 mbar oxygen pressure, increase in the background oxygen pressure makes the mean free path shorter and it results in the enhanced scattering of the ejected particles in the laser produced plasma by the oxygen gas molecules. It may cause reduction in the rate of film formation and the deterioration of the intensity may be due to the prevention of deposition of the atoms from reaching the right position in the crystal lattice from the flux at higher oxygen pressure [25,26].
Fig.1. XRD patterns of laser ablated BaSnO$_3$ thin films deposited on quartz substrate for different oxygen pressures

The average size of the crystallites ($D_{hkl}$) of the films can be calculated using the Debye Scherer’s formula [27],
\[ D_{hkl} = \frac{k\lambda}{\beta_{hkl} \cos \theta_{hkl}} \] 

----- (1)

Table 1. Structural parameters of BaSnO$_3$ thin films prepared under different oxygen pressures

| sample | \( I_{(110)} \) | \( I_{(200)} \) | \( \beta(110) \) (degrees) | Crystalline size (nm) | Lattice constant (nm) | Lattice strain |
|--------|-----------------|-----------------|--------------------------|----------------------|----------------------|----------------|
| BS 0   | 1               | 1               | 0.262                    | 30                   | 47±5                 | 0.4103         | 0.00619±0.0134 |
| BS 0.01| 2.56            | 2.17            | 0.256                    | 32                   | 63±18                | 0.4104         | 0.00286±0.0126 |
| BS 0.015| 3.33           | 3.71            | 0.251                    | 33                   | 66±33                | 0.4105         | 0.00373±0.0131 |
| BS 0.02| 6.8             | 4.62            | 0.221                    | 34                   | 69±23                | 0.4105         | 0.00282±0.0203 |
| BS 0.025| 0.83           | 0.52            | 0.277                    | 31                   | 53±22                | 0.4106         | 0.00538±0.022  |
| BS 0.03| 0.82            | 0.298           | 0.301                    | 27                   | 40±5                 | 0.4104         | 0.00854±0.0505 |
Where $\lambda$ is the wavelength of the X-ray radiation, $k$ is a constant which can be taken as 0.9 for spherical crystallites and $\theta_{hkl}$ is the diffraction angle. The average size of the crystallites in the prepared films varies from 27 to 34 nm and is shown in Table 1. The lattice constants for the films were calculated and found to be in agreement with the reported value 4.101 Å [28-30].
The factors such as the internal stress, thermal stress, strain and lattice distortion in the films can broaden the XRD peaks and hence the crystallite size estimated using the Scherer formula can be smaller than the actual value of the crystalline size [31,32]. The Williamson-Hall relation can be used to study the effect of strain induced broadening of crystallite size in the full width at half maximum (FWHM) of XRD peaks [33] and it is shown in Fig.2c.

\[ \beta_{hkl} \cos \theta_{hkl} = \frac{k \lambda}{D'} + 2 \eta \sin \theta_{hkl} \]  

(2)

\( D' \) is the size of the strain free crystallite, and \( \eta \) is the lattice strain. Plots of \( \sin \theta \) vs. \( \beta \cos \theta \) are drawn for all the films and straight line plots are obtained. The y-intercept of the plot is a measure of \( k \lambda / D' \) from which the crystallite size
can be calculated and the strain can be estimated from the slope of the plot (Table 1). From the data shown in Table 1 it can be found that all the films prepared have considerable strain in them and there is strain induced broadening for the XRD peaks.

**XPS Analysis**

X-ray photoelectron spectroscopy (XPS) was used for the analysis of the chemical states and the composition of the elements present in the laser ablated BaSnO₃ films prepared under oxygen ambience. The presence of the peaks corresponding to Ba, Sn and O in the XPS spectra of the films suggest the formation of BaSnO₃ phase in the films. The XPS survey spectrum of a typical BaSnO₃ film deposited under oxygen ambience at ρO₂ = 0.02mbar is shown in Fig.3a. The Cls peak at 284.4 eV observed in the XPS spectrum of the film can be due to the carbon contamination of film surface in the presence of air. Fig.3b shows the core level spectra of Ba3d peaks of Ba. Manju et al reported Ba3d₃/2 and Ba3d₅/2 peaks at binding energy values 795.2 and 777.9 eV with a peak separation of 15.3 eV corresponding to Ba²⁺ state [34]. Singh et al reported the Ba3d spin orbit doublet peaks of Ba3d₃/2 and Ba3d₅/2 at the binding energies 794.8 and 779.4eV respectively with a peak separation of 15.4 eV [35]. In the present case, for the BSO film, the core level spectrum of Ba (Fig. 3b) yields two peaks at binding energy values 795.06 and 779.78 eV respectively with a peak difference of 15.28 eV. These peaks can be unambiguously assigned to spin orbit doublet peaks of Ba3d₃/2 and Ba3d₅/2 respectively. In all the films prepared at different oxygen pressure namely ρO₂ = 0.01, 0.02 and 0.03 mbar, the Ba3d₃/2 and Ba3d₅/2 peaks are observed at binding energy values of 795.32 and 780.04 eV respectively (with the peak difference of 15.28 eV). The observed peak positions suggest that Ba atom in the compound is in the Ba²⁺ state.

According to Kwoka et al., Sn can exist in three oxidation states Sn⁰, Sn²⁺ and Sn⁴⁺ corresponding to binding energies 485.0, 485.9 and 486.6 eV respectively [36]. The Fig.3c shows the core level spectra of Sn3d state of the different films. For the BS 0 film for the spin orbit doublets Sn3d₅/2 and Sn3d₃/2 peaks are observed at binding energy values of 486.04 and 494.49 eV with peak difference of 8.45 eV. For all the films prepared under different oxygen pressures namely ρO₂ = 0.01, 0.02 and 0.03 mbar (BS 0.01, BS 0.02 and BS 0.03), the Sn3d peaks are observed at binding energy values of 486.67 eV and 495.12 eV with the peak difference of 8.45 eV. The observed peak positions suggest that tin is in the Sn⁴⁺ state. The O 1s peak (Fig. 3d) is observed at 531.33 eV for the BS0 film and 531.17 eV for the films prepared under different oxygen pressures 0.01, 0.02 and 0.03 mbar. It is reported that the peak in the
range 530.1 to 532.0eV corresponds to the oxygen vacancy [37,38,39] This confirms the presence of oxygen vacancy or defects in the films.

Fig.3  XPS spectra of laser ablated BaSnO$_3$ film (a). XPS survey spectrum of BS 0.02 film, (b). the core level splitting of peaks of Ba, (c). the core level splitting of peaks of Sn and (d). the core level splitting of peaks of O.

Morphological analysis

The FESEM images of the films deposited without background oxygen ambience and with oxygen ambience are shown in Fig. 4. The effect of oxygen ambience on the surface morphology of the BaSnO$_3$ films can be clearly seen in the FESEM images. The surface morphology of the BS0 film (Fig. 4a) shows the random distribution of isolated clusters of different sizes which are formed due to the agglomeration of smaller grains. The FESEM images of films
prepared under the oxygen ambience (Figs. 4(b)-4(f)) show the uniform dense distribution of grains with good connectivity. This indicates that the films with better quality are formed in the presence of oxygen ambience.

Fig.4. FESEM images of laser ablated BaSnO$_3$ films deposited at different oxygen pressures, (a) film in oxygen free ambience, (b) film at $\rho O_2 = 0.01$ mbar, (c) film at $\rho O_2 = 0.015$ mbar, (d) film at $\rho O_2 = 0.02$ mbar, (e) film at $\rho O_2 = 0.025$ mbar and (f) film at $\rho O_2 = 0.03$ mbar
AFM analysis is useful for the study of the surface morphology, surface homogeneity and nanostructure of deposited films. The AFM images of two typical films namely BS 0.015 and BS 0.02 are shown in Figs. 5(a) and 5(b).
respectively. The root mean square (RMS) surface roughness of the BaSnO₃ films is obtained using WSXM 5.0 Develop 6.4 software and the values are shown in Table 2. It can be seen that RMS surface roughness decreases with increase in the value of background oxygen pressure. The bombarding species on the substrate and the film surface may vary with different oxygen pressures and hence the surface roughness can change with oxygen pressure [40]. This indicates the formation of the films with more uniform surface morphology, at higher oxygen pressures. The thickness of films is measured using the cross-sectional FESEM images. The cross sectional FESEM images of BS 0.015 and BS 0.02 are shown in Fig.5c and 5d. The non-uniform formation of film with less thickness can be seen in the FESEM image of BS0 film. But, the formation of films with uniform thickness can be seen in the cross sectional FESEM images of films deposited at background oxygen pressure 0.01 mbar and 0.02 mbar. Thus optimum value of background oxygen pressure is helpful in getting high quality uniform films. It is found that the thickness of the film increases systematically with increase in background oxygen pressures (Fig.5f). In the presence of background oxygen more number of scattering centers are available and they scatter the incoming flux so that The mobility and the diffusion of surface atoms will increase with increase in oxygen pressure. This may be the reason for the increase of film thickness with increase in oxygen pressure [41].

The elemental analyses are done using energy dispersive X-ray (EDX) spectra (Figures not shown). The presence of peaks corresponding to Ba, Sn and O indicate the formation of BaSnO₃ films.

Table 2. Optical and Morphological parameters of BaSnO₃ thin films for different oxygen pressures

| Sample code | Average transmittance (%) | Band gap (eV) | RMS roughness (nm) | Thickness (nm) |
|-------------|---------------------------|---------------|-------------------|----------------|
| BS 0        | 65                        | 2.91          | 13.46             | 108            |
| BS 0.01     | 81                        | 3.04          | 12.57             | 161            |
| BS 0.015    | 77                        | 3.05          | 12.08             | 170            |
| BS 0.02     | 85                        | 3.09          | 10.47             | 204            |
| BS 0.025    | 82                        | 3.09          | 10.46             | 244            |
| BS 0.03     | 84                        | 3.06          | 10.37             | 264            |
Optical analysis

The optical absorbance and transmittance spectra of the films deposited at different oxygen pressures in the wavelength range 190-2300 nm are shown in Figs. 6(a)-6(b) respectively. All the films have good transparency in the visible and near infrared range. The sharp fall of transmittance near the absorption edge observed for all the films indicate their good crystalline nature. The average transmittance of the films in the wavelength region 250-2300 nm are measured and the values are shown in Table 2. The BS 0 film present an average transmittance value of 65 % whereas the films prepared under oxygen ambience possess higher values of average transmittance. This indicates the films prepared under oxygen ambience are of superior quality compared to that prepared in oxygen free ambience.

The absorption coefficient \( \alpha \), can be calculated from the transmittance spectra using the following relation

\[
\alpha = \frac{1}{d} \ln \frac{1}{T} \quad - (3)
\]

where \( d \) is the thickness of the film and \( T \) is the transmittance of the film. The band gap energy \( E_g \) of the can be estimated from the following relation [32]

\[
\alpha(h\nu) = A(h\nu - E_g)^n \quad - (3)
\]

\( h \) is the Planck’s constant, constant \( A \) is the band edge sharpness which is related with the order in crystalline structure of deposited films, \( \nu \) is the frequency of radiation, and the exponent \( n \) is the exponent depending on the nature of band transition. The values of \( n \) for direct allowed and direct forbidden transitions will be \( \frac{1}{2} \) and \( 3/2 \) respectively. For the indirect allowed and indirect forbidden transitions \( n \) can take values 2 and 3 respectively [32,42,43]. The band gap is determined for all the films from the plot of \( (\alpha)^{1/n} \) versus photon energy \( h\nu \) on the x-axis by extrapolating the linear portion of the graph to the \( h\nu = 0 \). It is observed that for all the films, the best straight line is obtained for \( n = 2 \), which corresponds to indirect allowed transition. The estimated values of the optical energy gap for the films varies from 2.91 to 3.09eV(Table 2). The optical energy gap for the BS 0 film is 2.91 eV and its value is higher in films prepared under oxygen ambience. The optical energy gap increases systematically with increase in \( \rho O_2 \) up to 0.025 mbar oxygen pressure thereafter decreases slightly with increase in \( \rho O_2 \).
Fig. 6. UV-Photoluminescence spectra of laser ablated BaSnO$_3$ films for different oxygen pressures (a) absorbance spectra, (b) transmittance spectra, (c) Tauc plot and (d) Photoluminescence spectra.

The photoluminescence spectra of the BaSnO$_3$ films recorded for an excitation wavelength 350 nm are given in Fig 6(d). In BaSnO$_3$ films the emissions are caused by the transitions between energy levels of constituent elements barium and tin of the compound [44]. A violet emission around 400 nm can be expected in this compound due to direct transition between the energy levels of O$_2$p and Sn$_5$s [45]. Sumithra et al [46] observed blue–green emission around 445 nm and they have attributed this to the charge transfer transitions of SnO$_6$ octahedra. They also reported a green emission around 550 nm due to the defect centres, or can be due to inter site positions leading to the formation of oxygen vacancy [46]. In the PL spectrum of the BS 0 film presents three intense emissions at 408, 430, 580 and 625 nm. The violet emission at 408 nm can be due to $^{2}\text{P}_0 ^0\rightarrow 2\text{D}_{3/2}$ transitions of barium [47] or due to the transition between O$_2$p and Sn$_5$s levels [45]. The emission at 430 nm can be due to the charge transfer transitions in SnO$_6$ octahedron [46] or due to $^{3}\text{P}_0 ^0\rightarrow 3\text{P}_1$ transition of barium [47]. Soumitra et al., reported assigned this red emission in the 625 nm is due to the transition related to the oxygen interstitials and also due to lattice distortion along the preferred axis. The
weak yellow emission at 580 nm may be due to the transition of the defect centres. Various authors assigned different reasons for the origin of red emission around 625 nm. Upendra kumar et al reported the intense red emission in 625 nm region is due to the \( ^3\text{D}_3 - ^3\text{D}_0 \) transition of barium [44]. Soumitra et al., reported the intense red emission in 625 nm region is due to the \( ^3\text{D}_3 - ^3\text{D}_0 \) transition of barium [46]. The luminescence intensity in the spectrum shows that the intensity of the films deposited without oxygen pressure has high intensity in the blue emission region and the intensity decreases with increasing background pressure. The reduction in intensity of the films with oxygen ambience may be due to the reduction of defects or oxygen vacancies present in the sample.

**Conclusion**

The BaSnO₃ films at different background oxygen pressures were prepared by pulsed laser deposition method. The XRD patterns show polycrystalline nature with cubic phase for all films. The film deposited at the oxygen pressure 0.02 mbar has the highest intensity peak with least FWHM, indicating the highest crystallinity for the films deposited at 0.02 mbar. The particle size also increases with increase with increase in oxygen pressure and lattice constant are calculated. The deterioration of intensity at higher oxygen pressure is due to enhanced scattering of the ejected particles in the laser produced plasma by the oxygen gas molecules. The band gap energy of the films are calculated. Morphological analysis shows that the films deposited at 0.02 mbar oxygen pressure has a uniform distribution of grains with defined grain boundary. All films show PL emission in the visible region. XPS analysis reveals that the oxidation state of tin as Sn⁴⁺ and the existence of oxygen vacancy in the prepared BaSnO₃ films.

**Conflict of interest**

There is no conflict of interest in the study for any of the authors.

**References**

1) E. Moreira, J. M. Henriques, D. L. Azevedo, E. W. S. Caetano, V. N. Freire, U. L. Fulco, and E. L. Albuquerque, J. Appl. Phys, 112, 043703 (2012).
2) S Upadhyay, O. Parkash and D. Kumar, J. Phys. D: Appl. Phys. 37 1483, (2004).
3) J. Cerda, J. Arbiol, R. Diaz, G. Dezanneau, J.R. Morante, Mat. Lett. 56 131–136 (2002).
4) B. Ostrick, M. Fleischer, U. Lampe, H. Meixner, Sens. Actuators, B 44 601–606 (1997).
5) C. Doroftei, P. D. Popa, F. Iacomi, J. Of OptoAndAdv Mater 14, No. 3-4, 413 – 417(2012).
6) Q. Liu, J. Dai, Z.Liu, X. Zhang, G. Zhu and G. Ding, J. Phys. D: Appl. Phys. 43 455401 (2010).
7) R. M. Katiliute, P. Seibutas, M. Ivanov, R. Grigalaitis, A. Stanulis, J. Banys, And A. Kareiva, Ferroelectrics, 464:49–58, (2014).
8) W.J Lee, H. J. Kim, J. Kang, D. H Jang, T.H Kim, J. H Lee, and K. H Kim, Annu. Rev. Mater. Res. 47:391–423 (2017).
9) B. Liu, Q Liu, Y. Zhang, Z. Liu, Lei Geng, J. of Alloys and Compounds 680 343-349(2016).
S. A. Salehizadeh, H. M. Chenari, M. Shabani, H. A. Ahangar, R. Zamiri, A. Rebelo, J. Suresh Kumar, M. P. F. Graca and J. M. F. Ferreira, RSC Adv. 8, 2100 (2018).

M. Jawed Ansaree & S Upadhyay, Ionics 21:2825–2838 (2015).

Lampe U, Gerblinger J, Meixner H B 24–25: 657–660 (1995).

Wensheng Lu & Helmut Schmidt, J Mater Sci 42:10007–10013 (2007).

S. Upadhyay, Om Parkash, J. of Mater. Science Letters 16, 1330–1332 (1997).

Kutty TRN, Vivekanadan, R Mat Res Bull 22:1457 (1987).

H. J. Kim et al., Appl. Phys. Express. 5, 061102 (2012).

M. A. Pen, J. L. G. Fierro, Chem. Rev. 2001, 101, 1981 (2017).

M. Jawed Ansaree & S Upadhyay, Ionics 21:2825–2838 (2015).

Lampe U, Gerblinger J, Meixner H B 24–25: 657–660 (1995).

Wensheng Lu & Helmut Schmidt, J Mater Sci 42:10007–10013 (2007).

S. Upadhyay, Om Parkash, J. of Mater. Science Letters 16, 1330–1332 (1997).

Kutty TRN, Vivekanadan, R Mat Res Bull 22:1457 (1987).

H. J. Kim et al., Appl. Phys. Express. 5, 061102 (2012).

M. A. Pen, J. L. G. Fierro, Chem. Rev. 2001, 101, 1981 (2017).

M. Jawed Ansaree & S Upadhyay, Ionics 21:2825–2838 (2015).

Lampe U, Gerblinger J, Meixner H B 24–25: 657–660 (1995).

Wensheng Lu & Helmut Schmidt, J Mater Sci 42:10007–10013 (2007).

S. Upadhyay, Om Parkash, J. of Mater. Science Letters 16, 1330–1332 (1997).

Kutty TRN, Vivekanadan, R Mat Res Bull 22:1457 (1987).

H. J. Kim et al., Appl. Phys. Express. 5, 061102 (2012).

M. A. Pen, J. L. G. Fierro, Chem. Rev. 2001, 101, 1981 (2017).