Wide-angle infrared absorber based on negative index plasmonic metamaterial

Yoav Avitzour,1 Yaroslav A. Urzhumov,2 and Gennady Shvets1

1Department of Physics, The University of Texas at Austin, Austin, Texas 78712
2COMSOL, Inc., 10850 Wilshire Blvd. #800, Los Angeles, CA 90024

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A metamaterials-based approach to making a wide-angle absorber of infrared radiation is described. The technique is based on an anisotropic Perfectly Impedance Matched Negative Index Material (PIMNIM). It is shown analytically that a sub-wavelength in all three dimensions PIMNIM enables absorption of close to 100% for incidence angles up to 45deg to the normal. A specific implementation of such frequency-tunable PIMNIM based on plasmonic metamaterials is presented. Applications to infrared imaging and coherent thermal sources are described.

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I. INTRODUCTION

The emergence of a new field of electromagnetic metamaterials was brought about by the demand for materials with exotic properties unattainable in nature. One such property is a negative refractive index that requires both the effective dielectric permittivity εeff and magnetic permeability µeff to be negative1,2. Applications of negative index metamaterials (NIMs) include “perfect” lenses, sub-wavelength transmission lines and resonators, miniature antennas, among others3,4,5.

The results reported in this paper pertain to two other recently emerged applications of metamaterials. The first one is wavelength-selective infrared and terahertz detection, important for thermal imaging, night vision, and nondestructive detection. Wide-angle power absorption efficiency is desirable for miniaturizing photodetectors or microbolometers down to the wavelength size. The second application is the development of coherent thermal emitters6,7,8 for spectroscopic and thermophotovoltaic (TPV) applications9,10,11,12. By virtue of the Kirchoff’s Law, emissivity of a thermal emitter approaches one is wavelength-selective infrared detection, important for thermal imaging, night vision, and terahertz. The emitted radiation absorption by a semi-infinite slab of a lossy metamaterial is a Negative Index Material (NIM)1,2,3,4 if the absorber material’s permittivity ε and permeability µ are matched to that of vacuum, i.e., when µzz = εyy ≈ 1, we find that A(λ0) ≈ 1. The efficiency is desirable for miniaturizing photodetectors or microbolometers down to the wavelength size. The emitted radiation of each detector is focusable to wavelength-sized spots using the far-field by large numerical aperture optics. With these applications in mind, we describe in this paper a ultra-thin metamaterials-based wide-angle absorber of near-infrared radiation. The design of this perfect absorber is inspired by a Perfectly Impedance Matched Negative Index Metamaterial (PIMNIM) constructed from plasmonic wires.

II. THEORETICAL BACKGROUND

As a background, we consider a simple problem of radiation absorption by a semi-infinite slab of a lossy metamaterial with engineered dielectric permittivity and magnetic permeability tensors ε and µ. Radiation is assumed to be incident in the x–z plane at an angle θ with respect to the vacuum–material interface normal ez. We further assume s–polarization of the incident wave as shown in the inset of Fig. 1 i.e., the only non-vanishing field components are Ey, Hz, Hx. Assuming that both ε and µ are diagonal tensors, their only relevant components are εyy, µxx, and µzz. For our semi-infinite slab (assuming that the metamaterial’s thickness Lx is sufficient to absorb all transmitted radiation), absorptivity A is limited only by reflection . A straightforward calculation yields the reflection r and absorption A = 1 − |r|2 coefficients:

\[
A = 1 - \left( \frac{\cos \theta - \sqrt{\epsilon_{yy}/\mu_{zz} - \sin^2 \theta/(\mu_{zz}\epsilon_{xx})}}{\cos \theta + \sqrt{\epsilon_{yy}/\mu_{zz} - \sin^2 \theta/(\mu_{zz}\epsilon_{xx})}} \right)^2, \tag{1}
\]

where all material parameters are, in general, wavelength-dependent.

Studying Eq. (1), we note that total absorption (A(λ0) ≈ 1) at a specific wavelength λ0 is achievable at normal incidence (θ = 0) if the absorber material’s impedance η = \sqrt{\mu_{zz}/\epsilon_{yy}} is matched to that of vacuum, i.e., when µzz(λ0) = εyy(λ0). Using impedance matching to achieve total absorption is a very well known approach in microwave engineering13,14. More surprising is another prediction of Eq. (1): that nearly total absorption can be achieved over a very broad range of angles. Assuming that µzz(λ0) = 1 and |µzz| = |εyy| ≫ 1, we find that A(λ0) ≈ 1 − tan^4 θ/2. For a 90deg full-angle (θmax = π/4) we find that 0.97 < A < 1 for 0 < θ < θmax. Even more remarkably, a similarly high broad-angle absorption is predicted even for much modest values of µzz and εyy. For example, if the absorbing medium is a Negative Index Material (NIM)15,16 with εyy = µzz = −1 + i, we find that 0.94 < A < 1. The angular dependence A(θ) for such perfectly impedance matched negative index material (PIMNIM) is shown in
Fig. 1 for the semi-infinite metamaterial slab, \( L_x = \infty \), as well as for the more general case of a finite slab, with \( L_x \approx \lambda_0 / 7 \). The two curves are barely distinguishable because of the high metamaterial loss.

![Image](image_url)

**Fig. 1:** (Color online) Angular dependence of the absorption coefficient for the idealized structure with \( \mu_{xx} = \epsilon_{yy} = -1 + i \), \( \mu_{xx} = 1 \), and \( L_x = \infty \) (dashed line), \( L_x = 200 \) nm (dashed-dotted line), barely distinguishable from the \( L_x = \infty \) line. Solid line: same for the specific implementation of the metamaterial absorber shown in Fig. 2. A schematic of the idealized wide-angle metamaterial absorber is given in the inset, S-polarized incident wave is assumed.

The main implication of these results is that a sub-wavelength slab of *almost any* impedance-matched metamaterial acts as a wide-angle absorber. The challenge now is to design such a metamaterial, especially in the technologically important infrared part of the spectrum. With this goal in mind, the rest of this paper is organized as follows. In section III we present one such possible metamaterial design based on two layers of alternating long and short plasmonic nanoantennas (Fig. 2), and demonstrate that the unit cell of this PIMNIM can be made highly sub-wavelength in near-infrared. The sub-wavelength requirement is critical for achieving wide-angle absorption: recent numerical simulations have found that \( A(\theta) \) rapidly drops with \( \theta \) when the unit cell is too large. Then, in section IV, we demonstrate that a single layer of PIMNIM enables optical absorption \( 0.7 < A(\theta) < 0.9 \) for a 90 deg full-angle scan, and that its small-angle absorption coefficient exactly agrees with the prediction of Eq. (1). Finally, in section V we show that extremely high optical intensity enhancements (at least \( \times 10^3 \)) can be produced at the absorption peak.

### III. PIMNIM DESIGN

The geometry of the unit cell of the proposed PIMNIM is shown in Fig. 2. Each unit cell consists of two parallel layers separated by the distance \( l_x \). Each layer consists of a cut-wire of width \( w_z \) and length \( w_y \) surrounded by the continuous in the \( y \)-direction wires of width \( d_z \). To enable planar fabrication, the heights \( t_x \) of the cut and continuous wires are assumed to be the same. Metal wires are assumed to be embedded inside a dielectric with \( \epsilon_d = 2.25 \). An s-polarized electromagnetic wave \( \vec{E} \parallel \hat{x}_y \) excites the electric and magnetic responses of the unit cell. The magnetic response is caused by counter-propagating currents flowing through the adjacent cut-wires. Both cut and continuous wires contribute to the electric response. A similar structure has been analyzed in the microwave part of the spectrum, where it is not sub-wavelength. Plasmonic effects, i.e. taking into account that metals have a finite dielectric permittivity \( \epsilon_m \equiv \epsilon_m' + i\epsilon_m'' \) with \( \epsilon_m' < 0 \), are necessary to miniaturize the unit cell.

**Fig. 2:** (Color online) Schematic of the PIMNIM structure. Unit cell for electromagnetic and electrostatic simulations is inside the dashed rectangles.

Electromagnetic resonances in effective permittivity \( \epsilon_{eff} \) and magnetic permeability \( \mu_{eff} \) of plasmonic composites are unambiguously related to the electrostatic surface plasmon resonances of the appropriate symmetry (electric dipole and magnetic dipole, correspondingly). Therefore, the first step in investigating the suitability of a particular unit cell geometry for a plasmonic NIM is to identify the frequencies of its electrostatic resonances. Specifically, magnetic activity (including negative index behavior) has been shown to exist only for \( \lambda > \lambda_{res} \), where \( \lambda_{res} \equiv 2\pi \epsilon / \omega_{es} \) is the vacuum wavelength corresponding to the frequency of the electrostatic resonance responsible for the magnetic activity. Such magnetically-active (MA) resonance of the PIMNIM structure has been calculated by solving the Poisson’s equation \( \nabla \cdot (\epsilon \nabla \phi) = 0 \) for the electrostatic potential \( \phi \), where \( \epsilon \) is a piecewise constant function equal to \( \epsilon_m \) inside the metal wires and \( \epsilon_d \) outside. Poisson’s equation can be solved as a generalized eigenvalue equa-
Potential distribution of the lowest-frequency MA resonance is shown in Fig. 3. Its magnetic nature can be deduced from the electric field loops formed between the two cut-wires. The above resonance occurs when the dielectric contrast between metal wires and the dielectric matrix is $\epsilon_m/\epsilon_d = -26.8$. Because the resonant dielectric contrast is determined from electrostatic calculations that contains no spatial scale, this value is determined by the geometric shape of the unit cell and its inclusions but not by the overall scale set by either of the periods $L_{x,y,z}$. Assuming gold as the plasmonic component and silica with $\epsilon_d = 2.25$ as an embedding dielectric, the electrostatic resonance occurs at $\lambda_{es} = 1.2 \mu m$, thereby setting the lower limit on the wavelength at which strong magnetic response is expected. The next, electrically active (EA) resonance corresponds to $\epsilon_m/\epsilon_d = -19$, is considerably blue-shifted from the MA resonance. Therefore, the PIMNIM structure depicted in Fig. 3 is promising as a NIM that has widely separated electric and magnetic resonances. Thus, its electromagnetic characteristics strongly resemble those of the original microwave NIMs. The important difference is that the proposed structure operates at the near-infrared frequencies and exploits plasmonic resonances to achieve sub-$\lambda_d$ cell size, where $\lambda_d$ is the wavelength inside the substrate.

It is instructive to compare the resonant frequency of the MA resonance of the PIMNIM structure with that of the traditional Double Fishnet (DF) structure. DF can be obtained from the PIMNIM by extending the cut-wires in z-direction until they merge with continuous wires running along the y-axis, thereby forming the second, orthogonal set of continuous wires in the z-direction. For example, the DF structure thus obtained from the PIMNIM structure shown in Fig. 3 is periodic in the $y-z$ plane with the periods of $L_y = 170 \text{ nm}$ and $L_z = 320 \text{ nm}$ containing two sets of intersecting metal strips in $y$ and $z$ direction, with the corresponding widths of $d_y = 80 \text{ nm}$ and $w_y = 200 \text{ nm}$. The MA plasmonic resonance is found at $\epsilon_m/\epsilon_d = -8.3$ which corresponds to $\lambda_{es} = 0.73 \mu m$ for gold. Of course, the physical sizes can be scaled from the above dimensions by an arbitrary factor because there is no physical scale in electrostatics. This simulation illustrates the challenge of making a strongly sub-$\lambda_d$ DF structure. For the unit cell to be sub-$\lambda_d$, the operational wavelength must be reasonably close to $\lambda_{es}$. Therefore, making a unit cell with the largest dimension of $\lambda/4$ requires $L_y < 200 \text{ nm}$. Such small DF structures have never been fabricated to date, which explains why sub-$\lambda_d$ DF-based NIM have never been produced.

**IV. SIMULATION RESULTS**

Full electromagnetic finite elements frequency domain (FEFD) simulations of light transmission/reflection through the PIMNIM structure shown in Fig. 2 (approximate physical parameters indicated in Fig. 3 and further optimized to achieve perfect impedance matching) were carried out using the COMSOL commercial package. Drude model for $\epsilon_m(\omega)$ of gold, with $\omega_p = 1.367 \times 10^{16} \text{ rad/sec}$ and $\gamma = 4 \times 10^{12} \text{ rad/sec}$ (taken from [26]), and $\epsilon_d = 2.25$ were used in the simulations. Reflection minimization was carried out using standard non-linear multivariable optimization to fine-tune the PIMNIM parameters and adjust the perfect-matching wavelength $\lambda_0$. The normal incidence ($\theta = 0$) results are shown in Fig. 4(a): zero reflection is achieved (by design) at $\lambda_0 = 1.5 \mu m$, with the single PIMNIM layer absorption coefficient $A \approx 0.9$ (or $A \approx 0.99$ for two layers).

Effective parameters $\epsilon_{yy} \equiv \epsilon_{eff}(\lambda)$ and $\mu_{zz} \equiv \mu_{eff}(\lambda)$ of the PIMNIM structure were extracted from the complex transmission and reflection coefficients through a single PIMNIM layer using the standard procedure [27]. Because the negative index band is the lowest (in frequency) transmission band of the PIMNIM, there is no ambiguity in determining $\epsilon_{eff}$ and $\mu_{eff}$. At the impedance matching point $\lambda_0 = 1.5 \mu m$ it is found that $\epsilon_{eff} = -0.85 + 1.4i$. Although the so-called figure-of-merit (FOM=Re $n$/Im $n$) is less than unity, achieving high FOM is not necessary (or even desirable) for accomplishing total wide-angle absorption. More important is ensuring that the unit cell of the PIMNIM is very sub-wavelength: $L_{x,y,z} \ll \lambda_d$. This is indeed accomplished by the present PIMNIM design: $n_d L_y / \lambda_0 \approx 0.3$.

The peak absorption, $A_{max}(\theta) = 1 - |r|^2(\theta)$, of the PIMNIM is plotted in Fig. 4. For small $\theta$ there is an excellent agreement between $A_{max}$ and Eq. (1), indicating that the PIMNIM is accurately described as an effective medium. Even for $\theta = \pm 45 \text{ deg}$ absorptivity remains at.
about 70%. This implies that, in a narrow spectral interval around \( \lambda_0 \), emissivity of a heated PIMNIM is close to that of a blackbody over a broad range of angles, which would enable to manufacture wavelength-size coherent thermal emitters.

V. FIELD ENHANCEMENT

Physically, the absorption peak corresponds to strong electric field enhancement inside the PIMNIM structure. For example, for the PIMNIM described by Fig. 4 it is found that the maximum intensity enhancement near the corners and edges of the cut-wire exceeds \( |E_{\text{max}}|^2/|E_0|^2 > 2000 \), where \( E_0 \) is the amplitude of the incident wave, as can be seen in Fig. 5. We note that calculating the maximum field in the simulation domain can be imprecise, because mesh irregularities can lead to mesh-dependent spikes at the corners and edges of the structure. In order to calculate the field enhancement while avoiding such numerical artifacts, the results presented in Fig. 5 were calculated by taking the maximum of \( |E| \) interpolated on a rectangular grid of \( 1 \times 1 \times 1 \text{ nm} \) and then smoothed by nearest-neighbor smoothing. This method was found to be robust under variations in mesh structure and density. Interestingly, the field enhancement is maximized at the absorption peak, at \( \lambda = 1.5 \mu m \), yet shows no structure near the electric resonance, at \( \lambda = 1.38 \mu m \). This can be explained by the very strong reflection at the electric resonance, which prevents the incident fields from penetrating into structure.

In addition to the enhancement near the corners and edges of the structure, fairly large field enhancement exists in a considerable volume between the cut wires. Fig. 6 presents \( |E| \) at two \( y-z \) cross-sections, for different values of \( x \). In Fig. 6(a) the magnitude of the electric field \( |E| \) is plotted at \( x = 0 \) (in the middle of the PIMNIM). Intensity enhancement of \( |E|^2/|E_0|^2 \approx 200 \) is apparent in about two thirds of the region between the front and the back cut-wires. Fig. 6(b) presents the field closer the front cut-wire, at \( x = -8 \text{ nm} \). An average intensity enhancement of \( |E|^2/|E_0|^2 \approx 400 \) is apparent in roughly the same area. If photon-counting detectors are integrated into the PIMNIM structure, then such intensity enhancement translates into proportionally enhanced absorption efficiency. The consequence of the large field enhancement is the desirable ultra-thin dimension of the absorber, in contrast to earlier calculations\(^28\) that demonstrated that thicker (about one wavelength) perfect absorbers can be developed using plasmonic spheres.
VI. CONCLUSIONS

In conclusion, we have demonstrated that an impedance-matched negative index metamaterial can act as a wavelength-selective wide-angle absorber of infra-red radiation. A specific implementation of such frequency-tunable PIMNIM based on plasmonic metamaterials is presented. Applications of the PIMNIM include infrared imaging and coherent thermal sources. This work was supported by the ARO MURI W911NF-04-01-0203, AFOSR MURI FA9550-06-1-0279, and the NSF NIRT grant No. 0709323.

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