Reciprocity in charge transport describes the symmetric behaviour of current with voltage: the magnitude of the current generated by a positive voltage is the same as that generated by a negative voltage. Violating reciprocity, that is, non-reciprocal behaviour, forms the basis for numerous important electronic devices such as diodes, a.c./d.c. converters, photodetectors, transistors and so on. A well-known example of a non-reciprocal device is the p–n junction, formed by the interface of a p- and n-doped semiconductor, which shows an asymmetric current–voltage characteristic (IVC) without applied magnetic field (B) and is widely used in various semiconducting technologies including logic and computation.

So far, in the superconducting regime, field-induced non-reciprocal behaviour has been seen in non-centrosymmetric superconductors and two-dimensional electron gases, and multiple mechanisms for this phenomenon have been proposed. All require an applied magnetic field to achieve simultaneous breaking of inversion symmetry and time reversal symmetry so that with sufficient magnetic field strength, the critical current required to destroy the superconducting state in one direction can be different from that in the other direction. Non-reciprocal superconductivity without an applied magnetic field as not yet been observed either in a bulk superconductor or a superconducting junction, such as a Josephson junction (JJ).

In particular, the JJ is also known to show reciprocal behaviour, as can be seen by the Josephson relations. JJs are typically formed by two superconducting electrodes bridged by a non-superconducting barrier composed of classical materials (AlOx, Cu and so on). However, integrating new quantum materials as the barriers in JJs may lead to anomalous behaviour violating the Josephson relations (such as non-reciprocal behaviour) and is of fundamental and technological interest. Creating a superconducting analogue to the semiconducting diode, a field-free ‘Josephson diode’ (JD), allowing supercurrent in one direction but normal current in the other direction in the absence of an applied magnetic field, is an open challenge.

One approach is to look for non-reciprocal IVCs in superconducting systems that only break inversion but not time reversal symmetry. Recently, JJs with only inversion symmetry breaking were predicted to show non-reciprocal IVCs owing to the formation of a Mott-insulating deple- 
tion region at the interface of p-doped and n-doped superconductors. As JJs form the basis of a variety of important technologies including superconducting quantum interference devices, superconducting quantum computing bit (qubits) and rapid single flux quantum devices, the realization of the JD will have a substantial impact on superconducting electronics.

Here we report the realization of the field-free JD in an inversion symmetry breaking the van der Waals heterostructure of NbSe2/Nb3Br8/NbSe2. An asymmetric IVC, with sharp superconducting transitions, is seen at a zero magnetic field indicating its field-free JD behaviour. A superconducting half-wave rectification of a square-wave excitation is demonstrated with a large rectification ratio, ultralow switching current density and highly stable performance. Moreover, the \( \Delta I_c \) versus \( B \) presents a symmetric behaviour, unlike the antisymmetric response seen in previous systems, further proving its potential for superconducting quantum devices.
field-free nature. The field-free JD effect may lie in the asymmetric tunnelling of supercurrent across the tunnel barrier, and the inversion symmetry breaking of the junction is confirmed by antisymmetric second-harmonic peaks. The effect shown here may be extended to a variety of material systems and architectures, particularly using van der Waals materials, for the advancement of future superconducting electronics and quantum devices. Also, this study shows how the interplay between symmetry breaking, new tunnel barrier materials and superconducting electrodes can lead to new or anomalous Josephson phenomena in JJs, opening the door to a variety of fundamental and applied research directions.

Experimental data

We fabricated vertical JJs by sandwiching an Nb,Br₄ thin flake with thin flakes of NbSe₂ as shown in Fig. 1a. Nb,Br₄ is a van der Waals quantum material that crystallizes in the R₃m space group at low temperature and its unit cell is composed of six layers of Nb-Br edge-sharing octahedra with inversion centres lying between adjacent layers (Extended Data Fig. 1a). The inversion symmetry is preserved in even-layered Nb,Br₄ but broken in odd-layered Nb,Br₄. Additionally, Nb,Br₄ was shown to have a singlet magnetic ground state and moderate band gap like its sister material, Nb₃Cl₈. (see Methods and Extended Data Fig. 5 for magnetoresistance characterization of the Nb,Br₄ thin flakes). The NbSe₂/Nb,Br₄/NbSe₂ heterostructures were fabricated by the dry transfer method (Methods) and capped by a top hexagonal boron nitride (h-BN) to avoid degradation, as a typical device in the inset of Fig. 1b shows. Here NbSe₂ flakes with thickness about 28–37 nm are used as superconducting electrodes, rotated by arbitrary angles with respect to each other, breaking inversion symmetry of the JJ, and a thin flake of Nb,Br₄ with thickness of 2.3 nm (three layers) is used as the tunnel barrier.

When cooling the junction to 20 mK in zero field, a sudden drop to zero resistance at $T_c = 6.6$ K is observed, as shown in Extended Data Fig. 1b. Figure 1b investigates the voltage versus current ($V$–$I$) behaviour, measured by sweeping the d.c. current at 20 mK (see $V$–$I$ curves at different temperatures in Extended Data Fig. 1c). We emphasize that the measurement of the $V$–$I$ curves contains four branches; sweeping the current from zero to positive (0–p), from positive back to zero (p–0), from zero to negative (0–n) and from negative back to zero (n–0). These four branches can be measured in the above order (defined as the positive sweep), or in reverse order, 0–n, n–0, p–0 and p–0 (defined as the negative sweep). Figure 1b shows $V$–$I$ curves of the positive sweep and the negative sweep; both of them show hysteresis, as expected, arising from the capacitance induced different critical currents when breaking ($I_c$) compared with returning ($I_c$) to the superconducting state. Note that the properties measured here all come from the JJ itself as the maximum current applied here is much smaller than the critical current of thick NbSe₂ (a few milliamperes). The two curves lie precisely on top of each other meaning that the positive versus negative sweep order does not affect the $V$–$I$ response. However, the $I_c$ and $I_p$ in the positive current regime (labelled $I_p$ and $I_c$) can be compared with those in the negative current regime (labelled $I_c$ and $I_p$). For reciprocal transport, as is expected for a conventional JJ without magnetic field, there should be no difference between $I_p$ versus $I_c$, $I_p$, $I_c$, $I_c$ and $I_p$, $I_c$, $I_p$, $I_c$.

Figure 1c shows the absolute values of $I_p$ and $I_c$ as well as $I_p$ and $I_c$ of the positive sweep. Immediately evident is that $I_p$ and $I_c$ are not equal to their negative current counterparts, with a $\Delta I_p$ ($\Delta I_p = |I_p| - |I_p|$) roughly 0.5 μA and a smaller $\Delta I_c$ roughly 0.1 μA, indicating the non-reciprocal $V$–$I$ response in our JJ. It is important to note that these differences are intrinsic properties of the JJ rather than an extrinsic Joule heating effect (see discussion in the Methods). The different absolute values of $I_p$ and $I_c$ combined with the sharp superconducting transition indicate that when the applied current is in between $I_p$ and $I_c$, the junction would show superconducting behaviour with a negative current but normal conducting behaviour with a positive current.

On the basis of this idea, we demonstrated half-wave rectification at zero field, as shown in Fig. 2a. Because the $I_p$ and $I_c$ being 7.61 and 8.14 μA, respectively, we applied a square-wave excitation (Fig 2a, top panel) with an amplitude of 7.9 μA at a frequency of 0.1 Hz. As shown in the bottom panel of Fig 2a, the junction remains in the superconducting state with the negative current (purple area) while switching to the normal state during the positive current (white area). The measured voltage in the superconducting state is smaller than 4 $\times$ 10⁻³ V (our detection limit), whereas the normal state junction voltage is 1.6 mV, meaning the diode rectification ratio demonstrated here is at least around 10⁴. In addition, the switching current density and switching power are 2.2 $\times$ 10⁻³ A cm⁻² and 12.3 nW, respectively, which are two and four orders of magnitude smaller than in the bulk superconducting diode.

We further probed the durability of the JD and robustness of half-wave rectification, as shown in Fig. 2b by conducting 10,000 continuous cycles with an applied square-wave excitation with an amplitude of 7.9 μA at 0.5 Hz and zero field. We note that the ‘on’ and ‘off’ state voltages stayed constant during the 10,000 cycles, and the device remains stable after that, indicating the high stability of the JD. Importantly, the JD effect is not just an ultralow temperature property; according
to the asymmetric $V-I$ curves observed at 0.9 and 3.86 K (Extended Data Fig. 2), an ideal half-wave rectification was also realized at 0.9 K whereas the half-wave rectification at 3.86 K had some punch-through error, likely due to thermal fluctuations, as shown in Extended Data Fig. 2). An ideal half-wave rectification was also realized at 0.9 K whereas the half-wave rectification at 3.86 K had some punch-through error, likely due to thermal fluctuations, as shown in Extended Data Fig. 2).

**Fig. 2** | Half-wave rectification and durability test of the JD at 20 mK and zero field. a. The top panel shows the applied square-wave excitation with an amplitude of 7.9 μA (in between $I_{C+}$ and $I_{C−}$) and frequency of 0.1 Hz. The bottom shaded region denotes the superconducting state in which the voltage is zero during negative current bias and the white region denotes the normal state in which the voltage is $V_d$ during positive current bias. The red dotted line is the zero line. b. The top panel is the applied square-wave excitation with an amplitude of 7.9 μA and frequency of 0.5 Hz over 10,000 cycles. The bottom panel is the coincidently measured junction voltage over those cycles, showing no degradation in the response. The red dashed lines are zero lines.

To further confirm the field-free nature of the JD effect, we measured $V-I$ curves containing 0–p and 0–n branches with an in-plane magnetic field (inset of Fig. 3a) swept from 40 to −40 mT with 0.5 mT steps. As before, the $I_{C+}$ and $I_{C−}$ was extracted and plotted in Fig. 3a ($V-I$ curves with magnetic fields shown in Extended Data Fig. 3). Both $I_{C+}$ and $I_{C−}$ decrease with increasing magnetic field, as expected for a JJ. Below 35 mT, Δ$I$ emerges, as the $I_{C−}$ increases faster than $I_{C+}$ as $B$ goes to zero. Note that $I_{C−}$ stays larger than $I_{C+}$ regardless of the direction of the magnetic field. Figure 3b shows the $\Delta I$ plotted as a function of $B$; the $\Delta I$ is sustained at zero field and its magnitude is independent of $B$ in the low-field region. The inset of Fig. 3b shows the variation of $\Delta I$ with decreasing $|I_{C−}|$, and that $|I_{C−}|$ can be up to 25% larger than $I_{C+}$. The field-dependent $\Delta I$ behaviour here differs from the previously mentioned field-induced superconducting diode effect in which the $\Delta I$ shows an antisymmetric field dependence. Also, in bulk superconductors with field-induced superconducting diode effects, $\Delta I$ versus $B$ has been sweep direction dependent, flipping sign, whereas in the field-free JD demonstrated here, $\Delta I$ versus $B$ shows no dependence on sweep direction (Extended Data Fig. 4).

To prove Josephson coupling of the junction, the $dV/dI$ mapping as a function of bias current and magnetic field was measured and is shown in Fig. 4a. As expected, the characteristic single slit Fraunhofer pattern from the Josephson effect is observed (see Methods and Extended Data Fig. 9 for further measurements and discussion). The critical current of the JJ decreases rapidly with increasing magnetic field, consistent with the $I_d$ decreasing in Fig. 3a. To further confirm the inversion symmetry breaking in the junction, the field-dependent first- and second-harmonic resistances were measured with an a.c. current of 5 μA and an amplitude of 5 μA and frequency of 7.919 Hz, as shown in Fig. 4b. The sharp drop of the first harmonic resistance in roughly ±30 mT shows the onset of superconductivity, whereas the two antisymmetric peaks appear in the second-harmonic resistance at the same fields indicating that inversion symmetry in the junction must be broken. There can be contributions to inversion symmetry breaking here; first is the intrinsic lack of inversion symmetry of the three-layer Nb$_3$Br$_8$, and second is the geometry of the vertical junction in which the top and bottom NbSe$_2$ electrodes are arbitrarily rotated relative to each other and the Nb$_3$Br$_8$ layer. Hence we also fabricated and measured four-layer Nb$_3$Br$_8$ devices, which show the same symmetric $\Delta I$ versus $B$ behaviour and field-free JD effect as the three-layer Nb$_3$Br$_8$ JJ as well as antisymmetric peaks in the second-harmonic resistance measurements (see details in the Methods and Extended Data Fig. 6). These results indicate that the asymmetric NbSe$_2$/Nb$_3$Br$_8$ interfaces in the junction play an essential role for inversion symmetry breaking and the field-free JD effect.

**Fig. 3** | Magnetic field dependence of $I_d$ and $\Delta I$. a. $I_d$ (orange dots) and $|I_{C−}|$ (green dots) obtained from the 0–p and 0–n branches of the positive sweep as a function of applied magnetic field. The field-free JD was swept from positive to negative. The diode effect ‘turns off’ by ±35 mT. Inset, a schematic of the measurement geometry; the orange and blue layers represent NbSe$_2$ and Nb$_3$Br$_8$, respectively. b. $\Delta I$ as a function of magnetic field. $\Delta I$ is roughly field invariant until 35 mT with the diode behaviour robust in zero field. Inset, $\Delta I$ as a function of $|I_{C−}|$ from modulation of magnetic field with the diode effect ‘turning off’ below 2.1 μA.
**Online content**

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Methods

Synthesis of Nb3Br8 crystal
Nb3Br8 crystals used in this work were grown through chemical vapour transport. Stoichiometric mixtures of niobium powder (Alfa, 99.99%) and NbBr5 (Strem, 99.9%) with a total mass of 1.5 g were ground together and added to a 14 mm inner diameter (with a 2 mm thick wall) fused silica tube in a glovebox and handled using standard air free techniques. Next, 20 mg of NH4Br was included as a transport agent. The tubes were then sealed air free at a length of approximately 30 cm, about 5 cm longer than the first two zones of a three-zone furnace. A three-zone furnace was used with a temperature gradient of 840, 785 and 795 °C with all but the last few centimetres of the tube between the first two zones. This discouraged the formation of large intergrown clumps of crystals at the end of the tube. The furnace was held at temperature for 3–5 days before being cooled to room temperature over 7 h.

Fabrications of the JJ device
We fabricated the NbSe2/Nb3Br8/NbSe2 JJ in a glove box with an inert environment to avoid oxidation and decay. The bottom NbSe2 flake was directly exfoliated from an NbSe2 single crystal (HQ grapheme) to a SiO2/Si wafer, then Nb3Br8 thin film and top NbSe2 flake were exfoliated and successively transferred layer by layer by using the polydimethylsiloxane assistant dry transfer method31. A h-BN (HQ graphene) layer was finally transferred on the top to protect the JJ from degradation in the atmosphere. Then Ti (3 nm)/Au (50 nm) electrodes were deposited on the top and bottom NbSe2 flakes for transport measurement.

Transport measurements
The transport properties of JJ device were measured in a BlueFors dilution refrigerator with a base temperature of 20 mK. A Keithley 6221 AC/DC Current Source Meter was used to inject both d.c. current and square-wave excitation, and the d.c. voltage was measured using a Keithley 2182A Nanovoltmeter. The a.c. measurements were performed by injecting a.c. current with a frequency of ω (7.919 Hz) using a Zurich MFLI amplifier, and the second-harmonic signals were obtained from the out-of-phase 2ω component of the a.c. voltage. A Fraunhofer pattern was obtained by measuring the differential resistance (dV/dI) versus curve under different magnetic fields using a lock-in amplifier (Zurich MFLI) with a d.c. bias (Keithley 2636).

Excluding the Joule heating effect
In Fig. 1b, we measured positive and negative sweep V–I curves with the same sweep rates. These two curves lie on top of each other, indicating that the V–I curves and associated Ic+ and Ic− are not dependent on the two sweep directions. As shown in Fig. 1c, the V–I has asymmetric Ic+ and Ic−. A non-zero dV/dI curve indicates that a NbSe2/NbSe2 junction can operate with a d.c. bias (Keithley 2636).

Magnetoresistance characterization of Nb3Br8 thin flake
The Nb3Br8 single crystal has previously been shown to have a phase transition from a high temperature ferromagnetic state to a magnetic singlet ground state at low temperatures (β phase)35. Here we use transport measurements to investigate the magnetic property of a Nb3Br8 thin flake. First, we measured sweep-up (magnetic field sweep from negative to positive) and sweep-down (magnetic field sweep from positive to negative) magnetoresistances of NbSe2/Nb3Br8/NbSe2 JJ (device 1, device in main text) with a temperature above Tc of NbSe2, superconducting electrodes. Extended Data Fig. 5a, b shows the magnetoresistances of device 1 at 10 K with in-plane and OOP magnetic fields up to 7 T, respectively. The overlapping sweep-up and sweep-down magnetoresistances curves reveal no hysteresis or reversion between these curves, indicating no ferromagnetic signal in the Nb3Br8 thin flakes36,37. Moreover, we also replaced NbSe2 superconducting electrodes with FLG to fabricate FLG/Nb3Br8/FLG devices and measured their magnetoresistances with an OOP magnetic field. Extended Data Fig. 5c shows the results of one of the FLG/Nb3Br8/FLG devices (device 6) measured at 2 K. Similar to the results in device 1, the magnetoresistances overlap well and no hysteresis or reversion signals can be observed in the curves, which is also consistent with the observation of the magnetic singlet ground state36. It is worth noting that the oscillations in FLG/Nb3Br8/FLG device are induced by the quantum oscillation of FLG.

Field-free JD effect in a NbSe2/four-layer Nb3Br8/NbSe2 JJ (device 3)
A NbSe2/Nb3Br8/NbSe2 JJ with a four-layer Nb3Br8 as the tunnel barrier (device 3) was fabricated using the same technique as device 1 (device in main text). We measured V–I curves at different magnetic fields at 20 mK, and extracted Ic+ and Ic− from 0–p and 0–n branches, respectively, and plotted them as a function of magnetic field as shown in Extended Data Fig. 6a. The Ic+ and Ic− show similar tendencies and behaviours to device 1. Extended Data Fig. 6b shows the field-dependent calculated ΔIc. The non-zero ΔIc can be clearly observed at zero field and low-field region, and it decreases to almost zero at relative higher fields, which is similar to the phenomenon in device 1. Therefore, the JJ with four-layer Nb3Br8 also has a field-free JD effect. To confirm the inversion symmetry breaking, we also performed the second-harmonic resistances (R2) measurements at 20 mK with an applied current of 4 μA, as shown in Extended Data Fig. 6c. The two asymmetric R2 peaks corresponding with the superconducting transition of R2 reveal the inversion symmetry breaking nature of the JJ9,15. Because the four-layer Nb3Br8 preserves inversion symmetry, as shown in Extended Data Fig. 1a, the inversion symmetry breaking of device 3 originates from the different interfaces of the top and bottom NbSe2/Nb3Br8.

Magnetic-field-induced superconducting diode effects in NbSe2/Nb3Br8 (devices 4 and 7) and NbSe2/FLG/NbSe2 (device 5) junctions
We also fabricated other junctions as control experiments to narrow the suspects of the field-free JD effect mechanism. First, we considered the influence of the superconducting electrode NbSe2, which is a type II and multi-gapped superconductor. A previous study reported non-reciprocal transport in atomically thin NbSe2 films that broke inversion symmetry, which disappeared in bulk NbSe2 crystal18. In all our devices, we used thick NbSe2 flakes (>20 nm) as superconducting electrodes, which have similar properties to bulk NbSe2. We also fabricated a NbSe2/NbSe2 junction (device 4) to explore its property. Extended Data Fig. 7a shows the V–I curve measured with the positive sweep (a sequence of 0–p, p–0, 0–n and n–0 branches) at 2 K and 0 T. As shown in Extended Data Fig. 7a, there are three transition steps in the V–I curve, which correspond with the critical currents of junction, bottom and top NbSe2 electrodes, respectively. There is no obvious difference between Ic+ and Ic− of the junction. The field dependence of critical currents was further measured; we extracted the Ic+ and Ic− of the junction from each V–I curve measured at different magnetic fields and plotted them as a function of magnetic field as shown in Extended Data Fig. 7b. The Ic+ and Ic− seemed to have a mirror symmetry with each other. ΔIc was also calculated and is shown in Extended Data Fig. 7c. The antisymmetric ΔIc curve indicates that a NbSe2/NbSe2 junction can only show a field-induced superconducting diode effect, rather than a field-free JD effect observed in NbSe2/Nb3Br8/NbSe2 junctions. Second, to further confirm the influence of barrier materials, we also fabricated a NbSe2/FLG/NbSe2 junction (device 5) as a control experiment. The Ic+ and Ic− obtained indeed are intrinsic to the device rather than induced extrinsically by the Joule heating effect.
and $I_c$ was extracted as before and shown in the inset of Extended Data Fig. 7d. $\Delta I_c$ was also calculated and plotted as a function of magnetic field in Extended Data Fig. 7d. Similar to the results from the NbSe$_2$/NbSe$_2$ junction device (device 4), the NbSe$_2$/FLG/NbSe$_2$ junction also shows the field-induced superconducting diode effect.

Previous studies have shown that a self-field effect in tunnel junctions with a large critical current density can result in a skewed Fraunhofer pattern, which may contribute to the field-induced superconducting diode effect in these junctions. We also fabricated a NbSe$_2$/NbSe$_2$ control device (device 7) with low critical current ($I_c \approx 120 \mu$A) and a correspondingly lower critical current density ($j_c \approx 18.5 \mu m^{-2}$) than device 4. $j_c \approx 66 \mu m^{-2}$. Extended Data Fig. 8 shows the $I_c$ and corresponding $\Delta I_c$ as a function of applied magnetic field. The $I_c$ and $\Delta I_c$ versus $B$ have a mirror symmetry with each other, and the obvious antisymmetric feature of $\Delta I_c$ versus $B$ curve also confirms the field-induced superconducting diode effect in device 7, similar to device 4.

In addition, recent studies on planar Nb/WTe$_2$/Nb JJs and the indium arsenide two-dimensional electron gas system, in which the self-field effect does not exist, also reported a mirrored $I_c$ versus $B$ curves, similar to the results from these NbSe$_2$/NbSe$_2$ and NbSe$_2$/FLG/NbSe$_2$ junctions. So far, the field-induced superconducting diode effect has been proposed to arise from effects such as magnetochemical anisotropy, finite momentum superconductivity and non-reciprocal Landau critical momentum. More investigation is needed to understand the origin of field-induced diode effect in NbSe$_2$/NbSe$_2$ and NbSe$_2$/FLG/NbSe$_2$ junctions.

In summary, both NbSe$_2$/NbSe$_2$ and NbSe$_2$/FLG/NbSe$_2$ junctions show the field-induced superconducting diode effect, similar to the behaviour of the (Nb/V/Ta) superlattice. These observations in NbSe$_2$/NbSe$_2$ and NbSe$_2$/FLG/NbSe$_2$ junctions indicate that the field-free JJD effect does not originate from the superconducting property of NbSe$_2$. Also, even though the interfaces of these junctions also break the inversion symmetry, the field-free JJD effect does not occur. These results, along with the observations in NbSe$_2$/NbBr$_3$/NbSe$_2$ devices (devices 1 and 3), indicate that the field-free JJD effect originates from asymmetric Josephson tunnelling induced by the NbBr$_3$ barrier and the associated NbSe$_2$/NbBr$_3$ interfaces in the junction.

**Basic analysis of the JJ (device 1) and discussion of Fraunhofer patterns (devices 1 and 2)**

The large hysteresis between 0–p (0–n) and p–0 (n–0) branches on the $V$–$I$ curves (Fig. 1b) indicates the JJ (device 1) lies in the underdamped region. On the basis of the critical current ($I_c$) and return current ($I_r$) of $V$–$I$ curve measured at 20 mK, the Stewart–McCumber parameter ($\delta = \frac{I_r}{I_c}$) of NbSe$_2$/NbBr$_3$/NbSe$_2$ JJ is 21.1, which is much larger than 1, confirming that the JJ is in underdamped region. On the basis of the equation $\delta = \frac{2}{\pi} \frac{R_c}{R_N}$ (where $C$ is the capacitance and $R_N$ is normal state resistance of junction), the junction capacitance calculated is $2.0 \times 10^{-14} F$, corresponding with specific capacitance ($C/A$, where $A$ is the junction area and is roughly 3.68 $\mu m^2$ in device 1) of 0.54 $\mu F$ cm$^{-2}$.

The Fraunhofer pattern in Fig. 4a is measured with a sweep-up magnetic field and Extended Data Fig. 9a shows the Fraunhofer pattern measured with a sweep-down magnetic field, which is similar to the sweep-up one. Both Fraunhofer patterns are not so periodic at higher field, which is a common case as seen in many other two-dimensional tunnel JJ devices, including NbSe$_2$/NbSe$_2$ junctions as well as NbSe$_2$/graphene/NbSe$_2$ junctions. In these vertical junctions with very thin barriers, the Fraunhofer pattern can be influenced by many effects, such as magnetic vortices in the superconducting electrodes and the self-field effect, which can make it difficult to get an ideal, highly periodic Fraunhofer pattern. In addition, from the theory, the ideal Fraunhofer pattern requires a uniform current profile (rectangular junction area shape and magnetic field parallel with either side of the rectangle). Therefore, the junction shape and magnetic field direction can also lead to a non-ideal Fraunhofer pattern, even in a high-quality junction with high-quality interfaces.

Here we also characterized more junction parameters by analysing the Fraunhofer pattern in Extended Data Fig. 9a at a low magnetic field. In a JJ, the London penetration depth ($\lambda$) of NbSe$_2$ can be obtained from $\phi_0 = \Delta B W (d + 2l)$, where $\phi_0$ is the magnetic flux quantum, $\Delta B$ is the period of the oscillation in the Fraunhofer, $W$ is the width of the junction perpendicular to the field direction and $d$ is the barrier thickness.

From the Fraunhofer pattern, the oscillation period is roughly $9.4 mT$ and with $W = 3.5 nm$ and $d = 2.3 nm$, $\lambda$ is calculated to be $3.85 nm$.

In addition, to further demonstrate the NbSe$_2$/NbBr$_3$/NbSe$_2$ heterostructure is indeed a JJ, we fabricated another NbSe$_2$/2-layer NbBr$_3$/NbSe$_2$ heterostructure device (device 2), and obtained the expected high-quality Fraunhofer pattern, as shown in Extended Data Fig. 9b. Moreover, the field-free JJD effect is also observed in this junction, as shown in Extended Data Fig. 9c and its inset. The $I_c$ versus $B$ curves shown in the inset of Extended Data Fig. 9c show a similar tendency and behaviour to device 1, and the extracted $\Delta I_c$ shows non-zero values at zero field and in the low-field region, and decreases to around zero at higher fields similar to device 1.

**Discussion on RCSJ mode of JD**

Here we discuss field-free JD behaviour based on the phenomenological RCSJ model of a JJ$^{17}$; the total current through the junction $I$ is determined by $I = I_c + I_{c+} + I_{c-}$. $I_c$ is the fluctuation current of the noise channel, $I_{c+}$ is the current of the capacitive channel, $I_{c-}$ is the extra resistive channel arising from finite temperature quasiparticle excitation and $I_r$ is the Josephson current. In d.c. measurement, the thermal noise current ($I_n$) dominates the fluctuation current, which can be calculated according to equation $I_n = (2e/\pi) \cdot h \cdot kBN$. Here, $h$ is Planck constant, $k_B$ is Boltzmann constant, $e$ is elementary charge and $T$ is the temperature. At 20 mK, $I_n$ is about 0.8 nA, far less than $I_{c-}$, and so will be disregarded as a cause of the $V$–$I$ asymmetry. In principle, the other three components could have asymmetric $V$–$I$ responses if inversion/time reversal symmetry is broken.

Below, we examine each of those possibilities in the context of our JJ.

Regarding the capacitive channel, as mentioned earlier, Misakia et al.$^{11}$ proposed that in an inversion symmetry broken JJ, an asymmetric charging energy of capacitance can lead to a $\Delta I_c$ without an applied magnetic field, which was observed in our JJ device (Fig. 1c). However, this model did not include an expectation of a $\Delta I_c$, which is one of the main requirements (the other being sharp transitions) for realizing half-wave rectification. As $I_c$ corresponds with the return to the superconducting state, applying a square-wave current excitation with an amplitude of $I_c + \Delta I_c/2$, for example, would not result in rectification as both ends of the wave lie in the superconducting regime.

Considering the extra resistive channel ($I_r$ = $V/R_N$) from quasiparticle excitations, its contribution to the diode effect would stem from asymmetry in $R_N$ and be significant if the resistance in the superconducting state was not zero. However, because the voltage in the superconducting state is zero in our device, the contribution of $I_r$ to the critical current can be neglected. Therefore, the critical current through the JJ is governed by the Josephson current ($I_J = I_c - I_{c+} + I_{c-}$), and the phase difference of two superconductors. The difference of positive and negative critical currents should originate from the asymmetric Josephson tunnelling in our JJ.

**The OAI phase of NbBr$_3$**

The crystal structure of NbBr$_3$ has the symmetries of the space group 166 ($R3m$), as shown in Extended Data Fig. 1a. in which Nb and Br atoms occupy the Wyckoff positions 6c and 18h, respectively. The electronic band structure and the band representation analysis of NbBr$_3$ have been studied, and the symmetry properties of its valence bands can be characterized by the multiplicities of irreducible representations at all the high-symmetry momenta, which is referred to as the symmetry data vector.

$$B = (m(F_G), m(F_G), m(F_G), m(F_G), m(F_G), m(F_G), ..., m(F_G))^T$$

(1)
where \( m(ρ) \) represents the multiplicity of the irreducible representations \( ρ \) formed by the occupied Bloch bands at the corresponding high-symmetry momentum. From topological quantum chemistry, the band representation, which are characterized by the above symmetry data vector, can be expressed as a linear combination of several elementary band representations (which are the band representations induced from atomic orbitals) with non-negative integer coefficients. Hence, \( \text{Nb}_3\text{Br}_8 \) is a topologically trivial insulator.

A topologically trivial insulator can be further classified as an atomic insulator or an OAI. If the band representation of a topological insulator cannot be induced from only the atomic orbitals at atom occupying sites, there must be some atomic orbitals (or Wannier functions) centred at the empty sites (that is, the Wyckoff positions) that are not occupied by any atoms) and this topologically trivial insulator is referred to as OAI. The OAI s can be indicated by the non-zero integer real space invariants at the empty sites and all the OAI s on the Topological materials database (https://www.topologicalquantumchemistry.org) have been exhausted by performing a high-throughput calculations, in which \( \text{Nb}_3\text{Br}_8 \) is identified as an OAI material indicated by a non-zero real space invariants at the empty Wyckoff position 3b, whose definition is

\[
\delta(b) = \frac{1}{2} \left[ -m(T_0F_0) + m(T_0) + m(F_0F_0) - m(T_0F_0) + m(T_0F_0) - m(T_0) \right].
\]

By substituting the symmetry data vector in equation (1) into equation (2), we have \( \delta(b) = -1 \), which means there must have a Wannier function centred at 3b. As the positions of 3b is not occupied by any atom, \( \delta(b) = -1 \) indicates that \( \text{Nb}_3\text{Br}_8 \) is an OAI and the Wyckoff position 3b is referred to as the obstructed Wannier charge centre.

To demonstrate the OAI phase of \( \text{Nb}_3\text{Br}_8 \), we have also calculated the electron density distribution in the lattice of \( \text{Nb}_3\text{Br}_8 \). As the results show in Extended Data Fig. 10, the electron charge density is symmetrically centred at the Wyckoff position 3b (the inversion centre between two \( \text{Nb}_3\text{Br}_8 \) layers), indicating the special separation of positive and negative charge centres along the c axis of the crystal. In the heterostructures of \( \text{NbSe}_2/\text{Nb}_3\text{Br}_8/\text{NbSe}_2 \), the asymmetric top and bottom interfaces break the inversion symmetry of the whole junction along c axis. The distribution of charge centres between layers is easy to be influenced by this symmetry breaking, which could result in the charge polarization in the whole junction along z direction and hence show an OOP polarization. As discussed in the main text, in analogy to previous theoretical investigations of the polarized system, we propose that a polarization in the \( \text{NbSe}_2/\text{Nb}_3\text{Br}_8/\text{NbSe}_2 \)Js may induce asymmetric Josephson tunnelling and lead to the field-free JD effect; further theoretical and experimental study is necessary to fully explain the mechanism.

Data availability

The data that support the findings of this study are available from the corresponding authors upon reasonable request.

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Author contributions

H.W. and M.N.A. conceived and designed the study. C.P. grew the samples. H.W. and Y.W. fabricated the devices. H.W., Y.W., P.K.S. and U.F. performed the transport measurements. H.W. and Y.W. carried out the data analysis. Y.X. provided theoretical support and discussion. Y.-J.Z. and S.S.P.P. provided facility and instrument support. T.M. and M.N.A. are the principal investigators. All authors contributed to the preparation of the manuscript.

Competing interests

The authors declare no competing interests.

Additional information

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Extended Data Fig. 1 | Temperature dependent resistance and \( V-I \) curves of the Josephson diode. a, \( \text{Nb}_3\text{Br}_8 \) crystal structure and inversion center locations. b, Resistance versus temperature measured using an a.c. current of 100 nA. Inset shows the enlarged plot near the superconducting transition of 6.6 K. c, \( V-I \) curves with positive sweep measured at different temperatures showing nonlinear behavior appearing below \( T_c \).
Extended Data Fig. 2 | V-I curves and half-wave rectification at different temperatures. a, Positive sweep V-I curve measured at 0.9 K. Both \( \Delta I_c \) and \( \Delta I_r \) are visible. b, Half-wave rectification measured at 0.9 K with an applied current of 7.6 \( \mu A \) at 0.1 Hz. The red dotted lines are the zero lines, showing that junction is in the superconducting state with negative current and switches to the normal state with positive current. c, Positive sweep V-I curve measured at 3.86 K. \( \Delta I_c \) is still visible, while the hysteresis is almost completely suppressed. d, Half-wave rectification measured at 3.86 K with an applied current of 5.14 \( \mu A \) at 0.1 Hz. The red dotted lines are the zero lines. Imperfect rectification is evident with some punch-through error, probably due to thermal fluctuation.
Extended Data Fig. 3 | V-I curves with 0-p and 0-n branches measured at different magnetic fields. Solid lines are 0-n branches (where $I_c-$ was extracted) and dotted lines are 0-p branches (where $I_c+$ was extracted) corresponding to Fig. 3. $\Delta I_c$ almost 'turns off' at 35 mT.
Extended Data Fig. 4 | Sweep-up magnetic field dependence of $I_c$ and $\Delta I_c$.

a, $I_c$ (orange dots) and $|I_c|$ (green dots) obtained from the 0-p and 0-n branches of the positive sweep as a function of applied magnetic field. The in-plane magnetic field was swept from negative to positive. 

b, $\Delta I_c$ as a function of magnetic field. Inset shows $\Delta I_c$ as a function of $|I_c|$ from modulation of magnetic field with the diode effect 'turning off' below 2.1 µA. These sweep-up results are nearly identical with the sweep-down results shown in Fig. 3.
Extended Data Fig. 5 | Magnetoresistance of NbSe₂/Nb₃Br₈/NbSe₂ and FLG/Nb₃Br₈/FLG heterostructures. a, Sweep-up and sweep-down magnetoresistances of NbSe₂/Nb₃Br₈/NbSe₂ (device 1) with OOP magnetic field at 10 K (above $T_c$ of NbSe₂). b, Sweep-up and sweep-down magnetoresistances of device 1 with in-plane magnetic field at 10 K. c, Sweep-up and sweep-down magnetoresistance of FLG/Nb₃Br₈/FLG (device 6) with OOP magnetic field at 2 K.
Extended Data Fig. 6 | Field-free Josephson diode effect of a NbSe$_2$/four-layer Nb$_3$Br$_8$/NbSe$_2$ device. a, $I_{c+}$ and $|I_{c-}|$ as a function of magnetic field at 20 mK in NbSe$_2$/four-layer Nb$_3$Br$_8$/NbSe$_2$ junction (device 3). b, $\Delta I_c$ as a function of magnetic field of device 3. c, First harmonic ($R_\omega$) and second-harmonic resistances ($R_{2\omega}$) as a function of magnetic field measured at 20 mK with an applied current of 4 µA of device 3.
Extended Data Fig. 7 | Field-induced superconducting diode effect in NbSe₂/NbSe₂ and NbSe₂/FLG/NbSe₂ heterostructures. 

**a**, $V$-$I$ curve of NbSe₂/NbSe₂ junction (device 4) measured at 2 K and 0 T. **b**, $I_c^+$ and $|I_c^-|$ as a function of magnetic field of device 4. **c**, $\Delta I_c$ as a function of magnetic field of device 4. **d**, $\Delta I_c$ as a function of magnetic field of NbSe₂/FLG/NbSe₂ junction (device 5) measured at 2 K. Inset shows the corresponding $I_c^+$ and $|I_c^-|$. 

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**Note:** The figure shows experimental data for various characteristics of superconducting diodes, including voltage-current ($V$-$I$) curves and changes in critical current ($\Delta I_c$) as a function of magnetic field ($B$) at different temperatures and device configurations.
Extended Data Fig. 8 | Field-induced superconducting diode effect in NbSe$_2$/NbSe$_2$ heterostructures with a lower critical current. a, $I_c^+$ and $I_c^-$ as a function of applied magnetic field of NbSe$_2$/NbSe$_2$ junction (device 7). b, $\Delta I_c$ as a function of magnetic field of device 7.
Extended Data Fig. 9 | Fraunhofer patterns of NbSe$_2$/Nb$_3$Br$_8$/NbSe$_2$ heterostructures. a, Fraunhofer pattern of device 1 with sweep-down magnetic field. b, Fraunhofer pattern of device 2. c, $\Delta I_c$ as a function of applied magnetic field of device 2 measured at 20 mK, inset shows the corresponding $I_c$ and $|I_c|$. 
Extended Data Fig. 10 | Obstructed atomic insulator property of Nb$_3$Br$_8$.

a, charge density distribution in a Nb$_3$Br$_8$ unit cell, the yellow/blue lobes indicate the charge density. b, charge density distribution as a function of unit cell location. Note the charge density does not drop to near zero in every other van der Waals gap.