Computational and Experimental Approaches to Controlling Bacterial Microcompartment Assembly

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ABSTRACT: Bacterial microcompartments compartmentalize the enzymes that aid chemical and energy production in many bacterial species. They are postulated to help bacteria survive in hostile environments. Metabolic engineers are interested in repurposing these organelles for non-native functions. Here, we use computational, theoretical, and experimental approaches to determine mechanisms that effectively control microcompartment self-assembly. We find, via multiscale modeling and mutagenesis studies, the interactions responsible for the binding of hexamer-forming proteins in a model system, the propanediol utilization bacterial microcompartments from Salmonella enterica serovar Typhimurium LT2. We determine how the changes in the microcompartment hexamer protein preferred angles and interaction strengths can modify the assembled morphologies. We demonstrate that such altered strengths and angles are achieved via amino acid mutations. A thermodynamic model provides guidelines to design microcompartments of various morphologies. These findings yield insight in controlled protein assembly and provide principles for assembling microcompartments for biochemical or energy applications as nanoreactors.

INTRODUCTION

Compartmentalization of cell components enables a variety of functions and ensures that biochemical processes happen without interfering with one another. Proteins play a crucial role in compartmentalization. Examples of closed compartments formed by protein components include viral capsids, which protect enclosed nucleic acids, and bacterial microcompartments (MCPs), which aid in the breakdown of chemicals and energy production that allow bacteria to thrive in various environments. MCPs, which are polyhedral structures capable of controlling the transport of specific molecules, are found in a large variety of bacterial species. Due to their native ability to enhance chemical transformations within cells, synthetic biologists are working to repurpose MCPs for non-native, industrially relevant purposes. Therefore, determining and understanding assembly mechanisms of MCPs are relevant to industry, life, and physical sciences problems.

The envelopes of MCPs are assembled from proteins, much like some viral capsids. Viral capsids that form closed crystalline shells are usually smaller than MCPs and mostly made of isomeric subunits that adopt spherical shapes with icosahedral symmetry or have an icosahedral shape. In contrast, MCP shells are often irregular in shape, contain multiple types of proteins, and enclose specific enzymes as cargo. Three types of shell protein components form the envelope of 1,2-propanediol utilization (Pdu) MCPs, ethanolamine utilization (Eut) MCPs, and many other MCPs. These three MCP shell components are composed entirely of two different protein domains, the bacterial microcompartment (BMC) domain (which comprises the hexameric and trimeric proteins) and the bacterial microcompartment vertex (BMV) domain (which comprises the pentameric proteins). Six BMC-domain monomers form the flat, six-sided, hexameric structures (BMC-H); three tandem BMC-domain dimers form the flat, six-sided, pseudohexameric trimers (BMC-T), and five BMV-domain monomers form the flat, five-sided, pentameric structures (BMC-P) (see Scheme S1 in the SI). When hexagonal shape components (BMC-H and BMC-T) assemble with pentagonal components (BMC-P) into closed shells, the topological constraints of Euler’s polyhedron formula are obeyed. This formula states that 12 pentagonal components (BMC-P) are needed to form a topologically closed shell. The position of the 12 pentagonal components in a crystalline shell is determined by the

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"quasiequivalent" triangulations of an icosahedron, which are all possible regular triangulations with 12 pentagonal (5-fold) disclinations.8 The 5-fold disclinations reside at the vertices of an inscribed icosahedron on a sphere,14,15 illustrating the ubiquitous icosahedral symmetry of spherical shells in viruses.8 These disclinations seed the buckling into an icosahedral shape when the shells are large, sufficiently inextensible, and/or highly bendable.16−18 Various models of the BMC-P, BMC-T, and BMC-H components’ assembly into icosahedral MCPs have been proposed in the literature (see Scheme S1 in the SI). However, these models are not always consistent with the observations, which attest that closed shells formed by many MCP systems, including the Pdu MCP system, have complex polyhedral geometries19−21 in which shell protein components are present at different ratios.22,23 Polyhedral MCPs have been predicted in multicomponent closed shells using elasticity theory,24−27 yet the conditions, including protein size, shape, and interactions, to achieve specific structures are unknown. In certain conditions, these MCP proteins assemble to produce extended cylindrical structures.19 A recent study demonstrated that open “icosahedral cages” can be formed in the absence of BMC-P proteins;23 these cages contain open spaces that enable BMC-P binding such that Euler’s polyhedron formula is still obeyed when the shell closes.

While genetic and crystallographic studies have increased the understanding of shell formation and MCP architecture, shell protein mutation studies have also provided information on specific shell protein functionality including structure and transport.24 Recent mutation studies that alter the charge of pore amino acids in Pdu MCPs, for example, showed how these mutations affect growth of bacteria harboring these altered MCPs on 1,2-propanediol as a sole carbon source, demonstrating their metabolic importance. Moreover, due to the high abundance of BMC-H proteins in the MCP shell,29 the mutations of hexamers that affect interhexamer interactions play a crucial role in determining the MCP assembly.30

Figure 1. All-atom (AA) molecular dynamics (MD) simulation results of a pair of PduA hexamers. (A) Equilibrium configuration. The arginine residue on each edge sticks out to form a hydrogen bond with the backbone of valine in the neighboring PduA hexamer. The two arginine residues are highlighted by increased bead size. Water and ions are hidden for visual clarity. (B) Illustration of the reference frame and angle between the two PduA hexamers. The angle $\theta$ is decomposed into the twisting angle $\theta_t$ and bending angle $\theta_b$ by projecting the normal of the neighboring protein ($\mathbf{n}_2$) onto the X−Z plane and the Y−Z plane, respectively (see Figure S1 for a different representation). (C) MD simulation values of $\theta$, $\theta_t$, and $\theta_b$. Negative $\theta_b$ indicates bending toward the concave side of PduA hexamers. (D) The potential of mean force between two PduA hexamers shows an attractive binding energy of about 11 kcal/mol.
Here, we develop a multiscale computational approach and a thermodynamic model combined with mutagenesis studies and experimental observations to analyze the assembly of native and mutated PduA MCP shell proteins. As an experimental model system, we study the PduA BMC-H protein of *Salmonella enterica* serovar Typhimurium LT2. By computationally studying specific interactions between PduA hexamers, we find the conditions for assembly into closed polyhedral or extended shapes and examine a specific mutation site on these hexamers that affects the assembly of MCPs into different morphologies. We use all-atom (AA) molecular dynamics (MD) simulations to determine the interaction energy and equilibrium bending and twisting angles of native and mutated hexamers and predict what mutations can lead to self-assembly of MCPs or extended shapes. Experimental results with mutated PduA hexamers qualitatively corroborate the simulation predictions. Our study reveals that electrostatic and hydrogen bonding interactions between the arginine and valine residues on the edge of PduA are key to the self-assembly of PduA hexamers and MCP formation, as the majority of the mutations to this residue negatively impact MCP formation. We use the structural information and interaction energy from AA MD simulations to construct a coarse-grained (CG) model for PduA self-assembly. Moreover, we determine what assembled geometries are accessible given a different stoichiometric ratio of the 3 major MCP component proteins. Finally, we use this CG framework in combination with a theoretical thermodynamic model to explore how modulating BMC-H/BMC-H and BMC-H/BMC-P interactions can permit access to an array of assembled morphologies. Simulation and theoretical findings agree with in *vivo* experiments of MCPs with encapsulated green fluorescent protein (GFP). These findings will help guide future studies that seek to repurpose MCPs with specific morphologies and shed light onto the rules governing MCP assembly in general.

### RESULTS

**Molecular Dynamic Simulations.** We first investigate the interaction between two native BMC-H subunits using AA MD simulations of the PduA crystal structure from *S. enterica* LT2 (PDB id: 3ngk). BMC-H proteins are found at high abundance in the MCP shell and play a crucial role in MCP assembly, raising the importance of interhexamer interactions in the molecular layer presumed to comprise the facets of the shell. The two PduA hexamers bind into a highly stable dimer (see Figure 1A and Video S1) as a result of the complementary shape and hydrogen bonding (see Table 1), in which the arginine at the 79th position (ARG79) binds to the backbone carbonyl oxygen of VAL25. The side chain of ARG79 fits in the pocket formed by the residues ASN29, VAL25, HIS81, and LYS26 (Figure 1A, inset). This result agrees with previous studies indicating that ARG79 plays a role in interhexamer binding. To better quantify the relative orientation of the two hexamers, we calculated the angle between the normal vector of the hexagonal planes (Figure 1B), which are determined by the carbon (the first carbon atom on the amino acid) on 3 residues (ARG79 for the wild type and the corresponding mutated residue for the mutants listed in Table 1) on alternate positions in the hexameric unit (see Figure 1A). The angle is projected onto the y–z plane and x–z plane to decompose into a bending angle \( \theta_b = 0.4 \pm 0.04 \)° (error bars are standard deviations) and a twisting angle \( \theta_t = 10 \pm 3 \)° at equilibrium (Figure 1C), with a mean total angle \( \theta = 11 \pm 4 \)°. (Histograms of the bending and twisting angles are shown in Figure S1.) This calculation suggests that in solution the PduA hexamer complex mostly adopts a near coplanar configuration, corresponding to a near-zero preferred curvature. In the assembled MCP, there is another configuration with the hexamers at a 30° angle relative to each other (see Figure S2). The latter bent configuration is likely due to collective interactions among other proteins and enzymes, as shown in our coarse-grained model described below, and therefore, it is not observed in simulations of two proteins.

We study the potential of mean force (PMF) of two PduA hexamers using umbrella sampling MD simulations (Figure 1D). The binding energy \((E_{\text{b,AA}})\) is estimated to be 11 ± 2 kcal/mol. This binding energy falls in the range of reported values of hydrogen bonds, further indicating that hydrogen bonds are a major contributor to PduA binding. The calculated binding energy and bending and twisting angles from atomistic simulations provide semiquantitative guidance for building larger-scale models.

Since AA simulations cannot include many proteins and cover the time scales required to assemble the proteins into different MCP morphologies, we use CG modeling to study the assembly of MCPs at larger time and length scales. CG simulations have been used to follow the assembly of simplified protein models into MCPs or viral capsids. CG models have also elucidated the initial faceting process and the formation of a complete MCP. However, these models have not investigated how interactions between multiple shell components result in various MCP morphologies such as cylinders, and the molecular origin of spontaneous curvature has not been explained. We hypothesize that the stoichiometric ratio of different components and the interaction strength between them control the morphology of assembly products.

We construct a CG model using the approximate shapes and charge distributions of the aforementioned PduA structure (PDB: 3ngk) for the BMC-H model, the PdB homologue from *Lactobacillus reuteri* (PDB: 4fax) for the BMC-T model, and the Pdn homologue GrpN from *Rhodopseudomonas rubrum* (PDB: 477a) for the BMC-P model. We validate this model by comparing assembled morphologies against known experi-

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**Table 1. Hydrogen Bond Analysis of Selected PduA Mutants (Listed as a Percentage of Time during the Simulation)**

| Mutant | Average H-bond Percentage |
|--------|---------------------------|
| ARG79-Side | 0.0% |
| ARG79-Main | 0.0% |
| ARG79-Side | 0.0% |
| ARG79-Main | 0.0% |

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The components of our CG model are shown in Figure 2A. In the CG simulations, the charges of the BMC-H, BMC-T, and BMC-P proteins are computed from the AA structures (see Figure S3); the mechanical properties such as the spontaneous curvature of BMC-H assemblies result from the inclined edge shape, that is, the inclination angle $\theta_i \approx 25^\circ$ obtained from the PDB (see CG simulation results using other inclination angles in Figures S4 and S5 in the SI). We first explore the various geometries with a different stoichiometric ratio of BMC-H and BMC-P. We find shells with icosahedral symmetry (that is, with 12 pentamers in the position of an inscribed icosahedron on a sphere) assembled with BMC-H:BMC-P = 5:3, which resemble the shape of in vitro MCPs from Haliangium ochraceum$^{19,35}$ (Figure 2B). When we reduce the content of BMC-P to BMC-H:BMC-P = 6:1, the addition of hexameric proteins allows the shell to assemble into the irregular shapes of Pdu MCPs reported in the literature,$^{12,20,41}$ illustrating how varying shell protein contents can control the morphology of MCP structures formed. When a small portion of BMC-H is replaced by BMC-T, similar structures are found (Figures 2C,D). However, when the content of BMC-T further increases, the asphericity increases (see Figure S6), and the BMC-T units seem to segregate, as predicted by elasticity theory of multicomponent shells.$^{25}$ A detailed structural analysis of assembled MCPs is available in Table S1. The impact of BMC-T on the assembly can be understood by the difference in bending angle and rigidity of the BMC-H and BMC-T (a bending angle analysis is given in the SI). We note that the BMC-H:BMC-P ratio used in these simulations is lower than that typically observed for Pdu MCPs in vivo, and the MCPs simulated are about 20–40 nm in diameter, compared to...
typical reported Pdu MCP diameters of 40–600 nm.21,41,44,45 These smaller MCPs equilibrate faster, allowing more efficient exploration of parameter space while still qualitatively reproducing the characteristic polyhedral shapes of MCPs observed in experiments.19,20 Simulations with BMC-H:BMC-P ratios of 8:1 produced bigger, more aspherical shells (Figure S7), similar to the shapes found in CG simulations of the MCP assembly.39

Long cylinders are obtained when only BMC-H proteins are in the assembly, as shown in Figure 2E. The shape is consistent with both in vivo and ex vivo experiments,34,46 in which overexpressing PduA proteins resulted in the assembly of long cylindrical tubes, and in vitro experiments in which the BMC-H from H. ochraceum spontaneously assembled into tubes in the absence of other shell proteins.19 The ability to reproduce vastly different shapes (cylinders and polyhedra) suggests that this simple CG model captures the key interactions of Pdu MCP proteins.

After validating our CG model by comparing the results with experimentally observed morphologies at different stoichiometries, we explore how the docking angle of BMC-H proteins impacts MCP assembly. AA MD simulations reveal (Table 2) that the docking angle of BMC-H can be tuned by mutating the arginine at the binding site. To this end, we use a combination of AA MD simulations, mutation experiments, and CG modeling to determine the role of hexamer–hexamer interactions on BMC-H assembly. AA MD simulations, the PduA conservation score (Figure S8), and prior work in the field11,32,34,35 indicate that the arginine at the 79th position (ARG79) is a key residue for controlling PduA interactions. Thus, we hypothesized that mutations to ARG79 of PduA could impact PduA and, subsequently, MCP assembly. We perform AA simulations of the hexamer–hexamer interface on 14 PduA mutants. Of these, 11 mutants retain some degree of hexamer–hexamer interfacing contact for the duration of the simulations, and their bending and twisting angles are shown in Figure 3A–C. Three mutants result in hexamers that entirely dissociate from each other, completely eliminating interfacial contacts (Table 2 and Figure S9 in the SI provide further details). The hexamer–hexamer dissociation arises when ARG79 is replaced by a negatively charged residue or hydrophilic residue with low isoelectric point (aspartic acid, glutamic acid, and glutamine), indicating that Coulombic interactions play an important role in stabilizing the hexamer–hexamer interface in PduA. The α helix of the opposite hexamer has a negative partial charge, which is thought to complement the native, positively charged ARG79 residue.36 Therefore, mutating ARG79 to the opposite charge results in total dissociation of the PduA.

**Mutagenesis Experiment Comparison with Simulations of the Hexamer Assembly.** We turned to an experimental approach to assess the assembly of PduA mutants and validate the above-described AA MD simulations. When PduA is overexpressed in E. coli cells, it naturally assembles into long cylinders that span the length of the cell body. The assembled protein structures form bundles within the cell cytoplasm, preventing cell division and leading to extended chains of cells (Figure S10). This allows for cell length to be used as a proxy for protein assembly and allows us to rapidly assess the assembly state of different PduA mutants. We hypothesized that mutants which formed stable dimers and had limited bending and twisting angles in AA MD simulations would be more likely to lead to the linked cell phenotype, indicative of protein self-assembly into cylinders. Indeed, the two PduA variants predicted to form hydrogen bonds (the native ARG79 and the R79K mutant, where “R79K” indicates that the arginine (R) at 79th position in the protein sequence is mutated to a lysine (K)) have among the lowest bending and twisting angles in simulations (Figure 3C, Table 2), and overexpression of both proteins leads to highly elongated cells (Figure 3E, Figure S10) compared to the negative control mutants (R79A and R79*, where “*” indicates a stop codon) (p < 0.0001). We also verified that the expression level of various PduA mutants did not correlate strongly with cell length, indicating that differences in assembly state are not likely to due to differences in expression (Figure S11). These results indicate successful assembly of wild-type (WT) PduA and PduA-R79K into cylinders. Surprisingly, we found that substitution of ARG79 with large hydrophobic residues (R79L) or aromatic residues (R79F, R79Y, R79W) also enabled some amount of assembly, as indicated by the presence of linked cells during overexpression (Figure 3E, Figure S10), although to a lesser degree than the native ARG79 (p < 0.001). While initially surprising, this assembly is predicted by the low bending and twisting angles calculated for these mutants from our simulations (Figure 3C). Free energy calculations for the binding energy of select mutants are also provided (R79K, R79W, and R79P; see Table S2), but these do not correlate as well with cell length as bending and twisting angles—while hydrophobic residues tryptophan (R79W) and proline (R79P) have similar binding energies, experiments indicate that they do not exhibit a similar self-assembly behavior. Overall, two residues that form hydrogen bonds (R79R and R79Y) and aromatic residues (R79F, R79Y, R79W) both confer small bending and twisting angles between hexamers. Together, these results imply that bending and twisting angles as determined by our AA MD simulations can provide insight into the assembly of PduA.

### Table 2. Summary of All-Atom Molecular Dynamics Simulation Results

| Mutation  | Hydrophobicity or Charge | Bending Deg | Twisting Deg | Assembly Structure | Binding Energy (kcal/mol) |
|-----------|--------------------------|-------------|--------------|--------------------|--------------------------|
| Native    | Positive Charge          | 0.4         | 10           | Y                  | −11 ± 2                 |
| R79K      | Positive Charge          | 24          | 7            | Y                  | −11 ± 2                 |
| R79D      | Negative Charge          | 24          | 7            | Y                  | −11 ± 2                 |
| R79E      | Negative Charge          | 24          | 7            | N                  | N/A                     |
| R79Q      | Hydrophobic (5.65)       | 24          | 7            | Y                  | −11 ± 2                 |
| R79N      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79S      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79T      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79Y      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79C      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79A      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79P      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79L      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79F      | Hydrophilic              | 24          | 7            | N                  | N/A                     |
| R79W      | Hydrophilic              | 24          | 7            | N                  | N/A                     |

**Note:** Negative bending angles indicate bending towards the outside of PduA. In the binding energy column, negative values correspond to attractive interactions.
The combination of experimental and AA simulation results described above suggests that PduA assembly is largely influenced by the inclination angle $\theta_i$ and twisting angle $\theta_t$ between hexamers. To understand what assembled morphologies may be accessed by these different PduA mutants, we performed CG MD simulations on PduA alone (no other proteins and no cargo) using the BMC-H model described above and documented how preferred inclination and twisting angle between hexamers impacts the morphology of PduA-only structures formed (Figure 3D). From the inclined geometry assumed in the CG model, the resulting bending angle between two hexamer planes is approximately $2\theta_i$ (see Figure 2A and Figure S4A). The CG model predicts that, at small $\theta_i$, the mutated PduA assemble into extended structures with low curvature, whereas at large $\theta_i$, the hexamers form structureless aggregates or nonextended pseudoclosed shells with high curvature (see Figure S12 for snapshots of assembled morphologies).
flat or curved sheets are expected to roll into cylinders or cones in the confinement of the cell. The predicted critical angles for self-assembly into extended structures from CG simulations are given by the red line in Figure 3C,D. We note that several nonassembling mutants have bending and twisting angles (calculated by AA simulations) below this line. This quantitative discrepancy is because AA simulations have detailed interatomic interactions omitted in the CG model that overestimate the CG critical bending and twisting angle. That is, agreement with the in vivo experiments is found if we take into account this shift by moving the red line to the lower left. Taken together, the AA and CG simulations indicate that pure PduA hexamers that form interfaces with small bending and twisting angles can assemble into extended structures. Mutations with big angles, despite having comparable attractive interaction strength, tend to form pseudoclosed shells, unstructured aggregates, and stacked layers (see Figure S12 in the SI).

Multicomponent Microcompartment Assembly: Thermodynamic Model Comparison with Coarse-Grained Simulations and Experiments. Extending our exploration of how shell protein interactions impact MCP assembly, we next used CG MD simulations and thermodynamic modeling to explore how hexamer–hexamer and pentamer–hexamer interaction strengths impact assembled morphologies. We hypothesized that modulating these interaction strength ratios would allow us to tune the morphologies of MCPs formed, enabling downstream engineering efforts. The thermodynamic model is constructed by numerically minimizing a free energy that includes protein–protein binding, protein–cargo interaction, elastic penalty of shell bending, and chemical potentials (see the Materials and Methods section), using parameters consistent with the CG model described in the previous section. Based on the hexamer–hexamer binding energies calculated in our R79 mutation simulations, we specified a range of accessible BMC-H/BMC-H ($\epsilon_{hh}$) and BMC-H/BMC-P ($\epsilon_{ph}$) interaction energies. We used a protein ratio of BMC-H:BMC-T:BMC-P:cargo = 4:2:1:2 for CG simulations (Figure 4A). For simplicity, the interaction parameters for BMC-T/BMC-T and BMC-T/BMC-H were set equal to the BMC-H/BMC-H interaction energy, because PduB trimers have conserved arginines in positions analogous to the arginines on PduA hexamers, and these are thought to be critical for binding (in Figure S6D, we discuss the effect of changing these values). The data shown in Figure 4A are for a BMC-H and BMC-T inclination angle of $\theta = 25^\circ$; however, data taken for an inclination angle of $15^\circ$ produce qualitatively similar morphologies and trends (Figure S4B), suggesting that the observed trends hold for a variety of inclination angles.

Figure 4. (A) Phase diagram of assembled shapes from CG MD simulations. The axes are the $\epsilon$ parameters defined in eq 1. The simulation box is constructed from a unit of 4 $\mu$C–H, 2 $\mu$C–T, 2 enzymes, and 1 $\mu$C–P, replicated 5 times in $x$, $y$, and $z$ directions. (B) Phase diagram from thermodynamics analysis (see parameters used in Table S5 in the SI). When the hexamers’ binding energy is weak, there is no assembly (red region). With stronger hexamer–hexamer interaction $\epsilon_{hh}$, MCP, cylindrical, and samosa shaped shells are formed, corresponding to the blue, yellow, and green regions, respectively. (C, D) Snapshots of CG simulations (the hexamers, pentamers, and cargo are in green, blue, and yellow, respectively; a red dot is marked on the center of BMC-T to distinguish them from BMC-H). (C) The “samosa” shape ($\epsilon_{hh} = 3.8, \epsilon_{ph} = 1.9$) is a quasiclosed surface without BMC-P proteins. They are different from MCPs in that they have no pentamers, and the vertices are sharp cones with a hole or defect at the tip. An example of a 4-fold defect is shown in the center. (D) Coexistence of cylinders and “samosas” at ($\epsilon_{hh} = 3.8, \epsilon_{ph} = 2.28$).
The different morphologies predicted by the CG simulation and by the thermodynamic theory as a function of pentamer–hexamer and hexamer–hexamer interactions are shown Figure 4A,B, respectively. The phase diagram constructed using CG MD simulations shows that closed MCP structures form when $\epsilon_{hh} > 0.7$, and $\epsilon_{hh}$ is between 3.6 and 4.2. This range of energy values corresponds to a hexamer–hexamer attraction energy between 4 and 6 kcal/mol, obtained using eq 1, which does not include the repulsion caused by screened electrostatics (adding this screened Coulomb energy contribution, the overall attractive energy is between 3.6 and 5.6 kcal/mol). When hexamer–hexamer interaction strength is increased beyond this MCP-forming range ($\epsilon_{hh} > 4.4$), the proteins bind irreversibly and assemble into gel-like structures. When hexamer–hexamer attractions are weakened ($\epsilon_{hh} < 3.4$), no assembled morphologies are observed. When $\epsilon_{ph}$ is small compared to $\epsilon_{hh}$ ($\epsilon_{hh} < 0.7$), assembled structures do not contain pentamers. These structures include cylinders and structures resembling “samosas” with 4-fold coordinated defective holes at the corners (Figure 4C). We note that these samosa-like structures are also observed when inextensible sheets are folded to form a quasiclosed surface, for example, in macroscopic inelastic membranes.47

The thermodynamic theory (described in the Materials and Methods section) predicts changes in morphology with changing pentamer–hexamer ($\epsilon_{ph}$) and hexamer–hexamer ($\epsilon_{hh}$) interaction energies. By comparing the formation energy of each tested morphology, we obtain the phase diagram shown in Figure 4B. In the region of $\epsilon_{hh} < 1.0 k_BT$ (in red), the nucleation barrier is large, and as a result, no shells could be formed. At a moderate value of $\epsilon_{hh}$, the shells can assemble into multiple shapes depending on the value of $\epsilon_{ph}$. When $\epsilon_{ph}$ is sufficiently high, a small, closed shell (MCP) (blue region) results. In the lower pentamer–hexamer interaction region, hexamers aggregate without pentamers, and in such a case, cargo is required to provide the curvature to close the shell. The competition between the cargo–hexamer interaction and the bending energy decides whether a cylindrical shape or a “samosa” shape is assembled. While the cylindrical shells can encapsulate dense cargo at a considerable bending energy cost, the “samosa” shell can encapsulate loose cargo at a smaller bending energy penalty. These results are in qualitative agreement with the CG simulations in Figure 4A, except that the CG simulations for large $\epsilon_{hh}$ values produce gels of connected sheets trapped in a local minimum. The lack of quantitative agreement can be explained by the fact that the thermodynamic model assumes equilibrium and constant reservoir concentrations during the assembly.

Many of the morphologies observed in our CG simulations (Figure 4) have been observed in the literature for different MCP systems assembled under various conditions (Figure 5).20,24,45,46 supporting our hypothesis that shell protein interaction strength plays an important role in determining MCP morphology. To provide experimental support for the observed effects of changing hexamer–hexamer binding energy on MCP assembly in our models, we used a green fluorescent protein (GFP) encapsulation assay to probe MCP formation in strains expressing the PduA mutants characterized above. In this assay, GFP is fused to an N-terminal signal sequence that is sufficient for targeting GFP to the lumen of MCPs.20,48 This results in bright, fluorescent puncta in the cell cytoplasm if MCPs are present (Figure 6A). However, if MCPs do not form properly, fluorescence will be observed instead at one or both poles of the cells; these are termed polar bodies (Figure 6A). In this way, it is possible to determine if a strain is capable of forming MCPs.

To this end, we constructed 19 mutant strains of Salmonella enterica serovar Typhimurium LT2 to test the effect of altering PduA assembly on overall MCP assembly. Alterations were made to the pdu operon as shown in Figure 6B. Because PduJ is capable of overcoming loss of PduA function,30 we have knocked out the pdu open reading frame from these strains. Thus, the WT and ΔJ strains serve as positive controls, as these both contain functional, WT PduA. Note that fluorescent puncta are visible throughout the cytoplasm of these strains, as expected when MCPs form (Figure 6A). The ΔΔJ strain, in which both the pduA and pduJ open reading frames have been knocked out, serves as a negative control. Only polar bodies are observed in this strain, indicating improper MCP assembly (Figure 6A). In the experimental strains, a point mutant of PduA is encoded at the pduA locus (Figure 6A), and puncta were counted in each cell to determine the effect of these mutants on MCP assembly. Notably, only strains containing PduA mutants capable of forming hydrogen bonds (R79R and R79K) are also capable of forming MCPs (Figure 6C). These strains had significantly more puncta than all other mutant strains ($p < 0.001$) (Figure 6C). This includes strains that were demonstrated to have lower, but detectable, levels of assembly (R79L, R79F, R79Y, R79W) (Figure 3E). This is likely due to the fact that these strains are not able to form hydrogen bonds with the adjacent hexamer and therefore have lower binding stabilities. To
confirm these results, we attempted to purify MCPs from a number of the PduA point mutant strains which did not show MCP assembly in the GFP encapsulation assay. However, we were unable to purify MCPs from these strains, as evidenced by the lack of the standard MCP banding pattern by sodium dodecyl sulfate gel electrophoresis (SDS-PAGE) (Figure S13).
Together, these results demonstrate that only PduA variants capable of forming hydrogen bonds enable MCP assembly. This validates the results shown in Figure 4A which indicates that reduction of $\varepsilon_{\text{ih}}$ disrupts MCP assembly and leads to no assembly.

**DISCUSSION AND CONCLUSION**

Combined multiscale MD simulations and *in vivo* mutation experiments show that hydrogen bonding and Coulomb interactions of arginine residues on the hexamer edges (ARG79) of PduA play a crucial role in native PduA self-assembly. In our experiments, mutations of ARG79 to most other residues negatively affect the assembly of Pdu proteins into tubes, cones, or polyhedral compartments. From the 14 mutations tested by AA MD simulations, we find that those that have smaller bending and twisting angles are shown experimentally to be more likely to assemble into extended structures.

We determine the shape of these extended structures (tubes or cones) and confirm that small twisting and bending angles promote their formation by performing coarse-grained (CG) MD simulations. The CG model also determines the conditions to form polyhedral and icosahedral shells. MCPs contain hexamers, pseudohexameric trimers, and pentamers, which are incorporated in our CG model. Our CG model includes twisting and bending by considering the shape, the position of the charges, and specific short-range interaction (including hydrogen bonding and hydrophobic interaction) of Pdu proteins in the form of BMC-H, BMC-T, and BMC-P structures. In agreement with our experimentally observed trends, the CG model predicts that small bending and twisting angles facilitate PduA self-assembly into extended cylinders. However, the agreement is only qualitative because the CG bending angle used in our CG simulations is overestimated; that is, it tends to be larger than the equilibrium bending angle obtained in the AA simulations, which include more detailed interactions.

According to our CG MD simulations, BMC-H proteins that maintain proper bending and twisting angles can form MCPs when BMC-P proteins and BMC-T proteins are present, and BMC-P interactions with BMC-H are strong enough. For BMC-P/BMC-H mixtures that can form MCPs, the ratio number of BMC-H to BMC-P determines the size and asphericity of assembled shells, with a higher ratio of BMC-H leading to larger, more elongated shells. We find that increasing the content of BMC-T increases the asphericity of assembled structures (Figure S6) because BMC-T prefers smaller bending angles. All crystalline closed shells should have 12 pentamers required by Euler's theorem. However, we see fluctuations in the number of pentamers in some MCPs. Moreover, in some assemblies, pentameric units are stuck in 6-coordinated sites. This may be due to local kinetic traps that are not overcome by annealing (see the Materials and Methods section). Only in a few cases, when BMC-H and BMC-P units coassemble without BMC-T units, are the 12 pentamers further away from each other forming spherical shells with icosahedral symmetry, which are too small and/or too rigid to buckle into icosahedral shapes.

Other morphologies including “samosas” and cylinders arise in our CG simulations and theoretical arguments by modifying the hexamer and pentamer interactions. These morphologies are robust in simulations using other CG bending angles (see Figures S4 and S5) and are also observed using a thermodynamic model that includes interactions, bending energies, and concentrations of the components via chemical potentials. Our experiments demonstrate that reducing the hexamer–hexamer interaction strength abolishes MCP assembly in GFP encapsulation assays, as predicted by our CG MD simulations and theory. Our work explains the necessary conditions for BMC domain proteins to assemble into closed MCPs, as well as cylinders, and suggests that it is possible to systematically modify the shape of the resulting assembled structures.

**MATERIALS AND METHODS**

**All-Atom (AA) Molecular Dynamics (MD) Simulations.** The PduA hexamers (PDB id: 3ngk) in the AA MD simulations are downloaded from the PDB and solvated in water containing 100 mM NaCl and 4 mM MgCl$_2$. Afterward, the system undergoes a short constant volume, temperature (NVT) equilibration of 100 ps with the backbone restrained. Then, restraints are released for one PduA hexamer, while the other still has a backbone restraint, and constant pressure, temperature (NPT) production simulations of 100 ns and umbrella sampling simulations are performed. For native PduA, an NPT extended simulation of 200 ns is performed; the PduA–PduA angle calculated is similar to 100 ns simulations. More details for AA MD simulations are in the SI.

**Coarse-Grained (CG) MD Simulations.** The BMC-H, BMC-T, and BMC-P models include short-range attractive interactions modeled by a Lennard-Jones–Gauss potential on a pair of attractive sites on the side of CG proteins (Figure 2A), given by

$$U_{\text{LJG}}(r) = \varepsilon_{ij}\left[\left(\frac{\sigma}{r}\right)^{12} - 2\left(\frac{\sigma}{r}\right)^{6} - 0.8\exp\left(-\frac{(r - r_0)^2}{2\sigma^2}\right)\right],$$

where $\sigma = 1.0$ nm is the bead size; $r_0 = 1.2$ nm roughly corresponding to the length of arginine side chain; $ij$ refers to the three species, $p$ for pentamer (BMC-P), $h$ for hexamer (BMC-H), and $t$ for trimer (BMC-T); and $\varepsilon_{ij}$ is the binding energy between species $i$ and $j$, where we choose the hexamer–hexamer interaction based on the AAMD simulation ($\varepsilon_{hh} = \varepsilon_{hh,\text{AA}} + \mu$, $\mu$ is the CG chemical potential). To reduce the enormous parameter space, the value of $\varepsilon_{ij}$ is set equal for the BMC-H and BMC-T proteins. The consequence of this assumption is discussed in the SI. $\varepsilon_{hh,\text{AA}}$ is unknown and left as a free parameter (see Table S2). In addition to the hydrogen bonding, the proteins’ charges (Figure 2A, right column) interact via a screened Coulomb potential with Debye length of 0.91 nm. The resulting total hexamer–hexamer interaction from two sites is then about $-7k_B T$ (or $-4.4$ kcal/mol) per edge. The spherical cargo is a generic model for enzymes encapsulated in MCPs; they are attracted to BMC-H and BMC-T proteins at six sites by a standard 12–6 Lennard-Jones (LJ) potential with $\varepsilon = 1.8$, resulting in a total cargo/shell protein binding energy of $\sim-7.5k_B T$. The energy scale for this LJ potential was selected such that the cargo can bind to BMC-H and BMC-T, but the MCP structures are not heavily deformed by the strength of the cargo/shell protein attraction. Note that the hexamer–hexamer interaction energy in the CG model is smaller than that calculated from AA simulations because the protein concentration in CG simulations is higher (See the appendix in the SI for details of MD simulations).
Plasmid and Strain Creation. All plasmids and strains used in this study are listed in Tables S3 and S4. For modifications to the Pdu operon in LT2, λ red recombineering was used as previously described (see the Extended Materials and Methods section in the SI).

PduA Self-Assembly Assay. The PduA self-assembly assay was carried out as previously described (see the Extended Materials and Methods section in the SI). Once cultures were grown to saturation, strains were imaged using phase contrast microscopy. All images were adjusted equally for brightness and contrast and were cropped to an area of 500 × 500 pixels. Measurements of cell length were done using the segmented line tool in ImageJ as described previously.30

Western Blot. Western blots were done on cell cultures expressing FLAG-tagged PduA variants (see the Extended Materials and Methods section in the SI).

GFP Encapsulation Assay. The GFP encapsulation assay was carried out to measure differences in MCP assembly between PduA variant strains as described in previous manuscripts (see the Extended Materials and Methods section in the SI).30

Thermodynamic Theory. In this section, we describe the thermodynamic model that explains the morphologies observed in CG MD simulations. Consider \(N_i\) protein monomers (corresponding to BMC-H hexamers in the MD simulations) and \(N_e\) enzyme cargo initially placed in the simulation box. When the system reaches equilibrium, the spontaneous angle determined by the cargo size (\(r_c\)) can be written as a sum of ideal free energy of each component and the interaction energy arising from shell bending energy, monomer–monomer binding energy, and the cargo–monomer interaction:

\[
G = N_b \kappa_b T \ln(\rho_b \epsilon_b) - 1 + U_0 + N_b \kappa_b T \ln(\rho_b \epsilon_b) - 1 \\
+ U + N_c \kappa_b T \ln(\rho_c \epsilon_c) - 1 \\
+ qN_c \kappa_b T (\epsilon_\text{bond} + \epsilon_\text{monomer}) + N_c \epsilon_c
\]

where \(\epsilon_0\) denotes the standard volumetric cell size (\(a^3\)), with \(a\) approximated by the monomer size; and \(\epsilon_\text{bond}\) and \(\epsilon_c\) represent the bending energy, binding energy per monomer, and cargo–monomer interaction, respectively. Note that we have assumed that the total internal energies for monomers (\(U_0 + U\)) and cargo (\(U_c\)) are constant.

Minimizing the free energy with respect to the number of free monomers, we obtain the equilibrium condition

\[
k_b T \ln(\rho_b \epsilon_b) = qk_b T \ln(\rho_c \epsilon_c) + N_b \kappa_b T \ln(\rho_c \epsilon_c) - 1 \\
- [q(\epsilon_\text{bond} + \epsilon_\text{monomer}) + N_c \epsilon_c]
\]

and the law of mass action for the particle density

\[
\rho_b = \rho_c e^{-\Delta G/k_b T}
\]

where \(\Delta G = q\epsilon - q\rho_b - N_c \epsilon_c\) is the formation energy with the excess energy given by \(\epsilon = \epsilon_\text{bond} + \epsilon_\text{monomer} + N_c \epsilon_c\), and the monomer and cargo chemical potentials are \(\mu_0 = k_b T \ln(\rho_c \epsilon_c)\) and \(\mu_c = k_b T \ln(\rho_c \epsilon_c)\), respectively.

To investigate the formation of the shell, we incorporate the line tension so that

\[
\Delta G = q\epsilon - q\mu_b - \sigma_b \sqrt{q(1-q)}
\]

where \(\mu = \mu_0 + N_c \mu_c\) and \(\sigma = \frac{4\kappa_b a}{\pi}\), \(Q\) is the monomer number of a complete shell, and \(A = \frac{3\sqrt{3}}{2}\) is the area factor of hexamer. Since the main contribution to the bending energy comes from the block geometry, it is reasonable to set the corresponding rigidity \(k_b\) as a piece-wise function, i.e., \(k_b = 0\) when no geometry overlaps, and \(k_b\) is infinite when overlap occurs. In such a case, only the corners of samosa shape and cylinder shape assemblies are subject to the bending penalty, with the local curvature defined by the cargo size \(r_c\). Thus, we write the formation energy for each morphology as

\[
\Delta G_{\text{shell}} = \left\{\begin{array}{ll}
q - 2\frac{a^2}{Q} + 2\frac{a^2}{Q} & (-\epsilon_0) + \sigma_b \sqrt{q(1-q)} - q\mu_0 \\
\frac{1}{4\pi a} & q\mu_b - q\mu_0
\end{array}\right.
\]

\[
\Delta G_{\text{cylinder}} = k_b 4\frac{a^2}{Q} (1 - \cos(\theta_c - \theta_b)) + \left\{\begin{array}{ll}
\frac{4}{a^2} & (-\epsilon_0) + \frac{4}{a^2} \cdot (-\epsilon_0) \\
\frac{4q}{a^2} - \frac{4q}{a^2} & (-\epsilon_0)
\end{array}\right.
\]

\[
\Delta G_{\text{cylinder}} = (6q - 2\sqrt{qA}) (-\epsilon_0) - q\mu_0
\]

where \(\theta_c\) is the angle between normal vectors of connected subunits, and \(\theta_b\) is the spontaneous angle determined by the intrinsic properties of the protein. (see Table S5 for the list of thermodynamic parameters).

## ASSOCIATED CONTENT

1. Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acscentsci.0c01699.

Additional schematics, figures, tables, details of MD simulation, and experimental methods (PDF)

Video S1: Constant pressure, temperature all-atom simulation of native 2 PduA hexamer proteins (pdb id 3ngk) in an aqueous solution with 100 mM NaCl and 4mM MgCl₂ (MPG)

Video S2: Coarse-grained system of 500 BMC-H, 250 BMC-T, 125 BMC-P, and 250 cargo quickly assembling into polyhedral MCPs, followed by slow relaxation of correcting malformed units such as flipped BMC-T or BMC-P with incorrect coordination number (MPG)

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Notes
The authors declare no competing financial interest.

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