Research Article

Heavy Scalar, Vector, and Axial-Vector Mesons in Hot and Dense Nuclear Medium

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In this work we shall investigate the mass modifications of scalar mesons ($D_0$, $B_0$), vector mesons ($D^*$, $B^*$), and axial-vector mesons ($D_1$, $B_1$) at finite density and temperature of the nuclear medium. The above mesons are modified in the nuclear medium through the modification of quark and gluon condensates. We will find the medium modification of quark and gluon condensates within chiral SU(3) model through the medium modification of scalar-isoscalar fields $\sigma$ and $\zeta$ at finite density and temperature. These medium modified quark and gluon condensates will further be used through QCD sum rules for the evaluation of in-medium properties of the above mentioned scalar, vector, and axial vector mesons. We will also discuss the effects of density and temperature of the nuclear medium on the scattering lengths of the above scalar, vector, and axial-vector mesons. The study of the medium modifications of the above mesons may be helpful for understanding their production rates in heavy-ion collision experiments. The results of present investigations of medium modifications of scalar, vector, and axial-vector mesons at finite density and temperature can be verified in the compressed baryonic matter (CBM) experiment of FAIR facility at GSI, Germany.

1. Introduction

The motive behind the heavy-ion collision experiments at different experimental facilities is to explore the different phases of QCD phase diagram. These experiments help us to understand the nuclear matter properties for different values of temperatures and densities. The hadronic matter produced in heavy-ion collisions may undergo different phase transitions, for example, the liquid-gas phase transition, the kaon condensation, the restoration of chiral symmetry, and maybe the formation of quark gluon plasma [1–4]. The compressed baryonic matter (CBM) experiment of the FAIR project at GSI, Germany, may explore the phase of hadronic matter at high baryon densities and moderate temperatures. These kinds of phases may exist in the compact astrophysical objects, for example, neutron stars. The property of restoration of chiral symmetry is closely related to the medium modifications of hadrons [4]. The medium modifications of Kaons, $D$ mesons, and light vector mesons had been studied using different theoretical approaches, for example, chiral model [5–13], QCD sum rules [14–19], and coupled channel approach [20–22]. Due to interactions the properties of hadrons in the medium are found to be different as compared to their free space properties.

The medium modifications of heavy scalar, vector, and axial-vector mesons at finite density and temperature of the medium had been studied very rarely [23–26]. In the present investigation we will study the mass modifications of heavy scalar mesons ($D_0$, $B_0$), vector mesons ($D^*$, $B^*$), and axial-vector mesons ($D_1$, $B_1$) at finite densities and temperatures. The study of in-medium properties of scalar, vector, and axial-vector mesons will be helpful to understand their experimental production rates. The medium modification of charmed mesons may modify the experimental production of ground state charmonium $J/\psi$ and the excited charmonium states ($\psi'$ and $\chi_c$). The charmonium, $J/\psi$, may be produced due to the decay of the higher charmonium states. However, the vacuum threshold value of heavy meson pairs lies above the vacuum mass of the excited charmonium states. Now if these heavy mesons get modified (undergo mass drop in the medium) then the excited charmonium states may decay to the open charmed meson pairs instead of decaying...
to the ground state charmonium. Thus to understand the production of charmonium states in heavy-ion collisions it is very necessary to study the medium modification of the heavy scalar, vector, and axial-vector mesons. The medium modifications of heavy vector mesons may also help us in understanding the dilepton spectra produced in heavy-ion collision experiments [27–32]. The dileptons are considered as interesting probe to study the evolution of matter produced in heavy ion collision experiments as they do not undergo strong interactions in the medium. In [33] the production of open charm and charmonium in hot hadronic medium had been investigated using the statistical hadronization model at SPS/FAIR energies. In this work it was observed that the medium modifications of charmed hadrons do not lead to appreciable changes in cross-section for D mesons production. This is because of large charm quark mass and different times scales for charm quark and charm hadron production. However, the charmonia yield is affected appreciably due to in-medium modifications.

The properties of scalar charm resonances $D_{s0}(2317)$ and $D_{s0}(2400)$ and hidden charm resonance, $X(3700)$, had been studied in [25] using coupled channel approach. In these studies the $D_{s0}(2317)$ and $X(3700)$ were found to undergo a width of about 100 and 200 MeV, respectively, at nuclear matter density. However, for the $D_{s0}(2400)$ mesons there was already large width of resonance in the free space and the medium effects were found to be weak as compared to $D_{s0}(2317)$ and $X(3700)$. In [18] the mass splitting of $D \rightarrow \bar{D}$ and $B \rightarrow \bar{B}$ mesons had been studied using the QCD sum rules in the cold nuclear matter and the calculated values of mass splitting at nuclear saturation density were 60 and 130 MeV, respectively. The Borel transformed QCD sum rules had also been used to study the properties of pseudoscalar $D$ mesons [19] and vector mesons, $\rho$, $\omega$, and $\phi$ [17]. The properties of the scalar mesons ($D_{s0}, B_{s0}$) in the cold nuclear matter using QCD sum rules have been investigated in [23]. The vectors mesons ($D^*, B^*$) and axial-vector mesons ($D_{1}, B_{1}$) had also been studied using QCD sum rules in cold nuclear matter in [24]. Note that in [23, 24] the properties of the meson were investigated at zero temperature and at normal nuclear matter density. However, in the present investigation we will find the in-medium masses of the scalar ($D_{s0}, B_{s0}$) and vector ($D^*, B^*$) and axial-vector ($D_{1}, B_{1}$) mesons at finite temperatures as well as at the densities greater than the nuclear saturation density.

In the present work to investigate the properties of scalar, vector, and axial-vector mesons we will use the QCD sum rules and chiral SU(3) model [5]. Within QCD sum rules, the in-medium properties of mesons are related to the in-medium properties of quark and gluon condensates. We will investigate the in-medium properties of quark and gluon condensates using the chiral SU(3) model. Using chiral SU(3) model we will find the values of quark and gluon condensates at finite values of temperatures and baryonic densities. These values of condensates will further be used to find the medium modification of mesons using QCD sum rules. The chiral SU(3) model along with QCD sum rules had been used in the literature to investigate the in-medium modification of the charmonium states $J/\psi$ and $\eta_c$ [34].

The present paper is organized as follows. In Section 2 we will give a brief review of chiral SU(3) model. Then in Section 3 we will discuss how we will evaluate the in-medium modifications of the scalar, vector, and axial-vector mesons within QCD sum rules and using the properties of quark and gluon condensates as evaluated in the chiral SU(3) model. In Section 4 we will discuss the results of the present investigation and finally in Section 5 we will give a brief summary of present work.

## 2. Chiral SU(3) Model

In this section we will briefly review the chiral SU(3) model used in the present investigation for the in-medium properties of heavy mesons. The chiral SU(3) model is based on the broken scale invariance and nonlinear realization of chiral symmetry [35–40]. The model involves the Lagrangian densities describing, for example, kinetic energy terms, baryon-meson interactions, self-interactions of scalar mesons, vector mesons, symmetry breaking terms, and also the scale invariance breaking logarithmic potential terms.

For the investigation of hadron properties at finite temperature and densities we use the mean field approximation. Under this approximation all the meson fields are treated as classical fields and only the scalar and the vector fields contribute to the baryon-meson interactions. From the interaction Lagrangian densities, using the mean-field approximation, we derive the equations of motions for the scalar fields $\sigma$ and $\zeta$ and the dilaton field, $\chi$, in isospin symmetric nuclear medium. We solve these coupled equations to obtain the density and temperature dependence of scalar fields $\sigma$ and $\zeta$ and the dilaton field, $\chi$, in isospin symmetric nuclear medium [12]. The concept of broken scale invariance leading to the trace anomaly in (massless) QCD, $\theta^\mu_\mu = \frac{8}{3} \pi G_\mu^\mu G^{\mu\nu}$, where $G_{\mu\nu}$ is the gluon field strength tensor of QCD, is simulated in the effective Lagrangian at tree level [41] through the introduction of the scale breaking terms [12]. Within chiral SU(3) model the scale breaking terms are written in terms of the dilaton field $\chi$ and also the scalar fields $\sigma$ and $\zeta$. From this we obtain the energy momentum tensor and this is compared with the energy momentum tensor of QCD which is written in terms of gluon condensates. In this way we extract the value of gluon condensates in terms of the scalar fields $\sigma$ and $\zeta$ and the dilaton field, $\chi$, and is given by the following equation [12]:

\[
\frac{\langle \alpha_s G_{\mu\nu} G^{\mu\nu} \rangle}{\pi} = \frac{8}{9} \left[ (1 - d) \chi^4 + \left( \frac{\chi}{\chi_0} \right)^2 \right] \times \left( \mu_{3\pi}^2 f_\pi^\sigma + \left( \sqrt{2} m_{\pi}^2 f_K - \frac{1}{\sqrt{2}} m_{\pi}^2 f_\pi \right) \zeta \right),
\]

where the value of parameter $d$ is 0.064 [10] and $m_{\pi}$ and $m_K$ denote the masses of pions and kaons and have values
139 and 498 MeV, respectively. $f_\pi$ and $f_K$ are the decay
constants having values 93.3 and 122 MeV, respectively. The
symbols $\sigma$, $\xi$, and $\chi$ denote the nonstrange scalar-isoscalar
field, strange scalar-isoscalar field, and the dilaton field,
respectively. $\chi_0$ denotes the value of the dilaton field in
vacuum. The vacuum values of $\sigma$, $\xi$, and $\chi$ are $\sim$93.3, $\sim$106.6,
and 409.8 MeV, respectively. Note that in above equation the
gluon condensate is written considering finite quark masses.
If we have massless QCD, then only first term written in
terms of dilaton field $\chi$ contributes to the gluon condensates.
Using the above equation we obtain the values of scalar gluon
condensates at different values of densities and temperatures
of the nuclear medium.

3. QCD Sum Rules for Scalar ($D_0$, $B_0$), Vector ($D^*$, $B^*$), and Axial-Vector ($D_1$, $B_1$) Mesons

In this section we will discuss the QCD sum rules [23, 24]
which will be used later along with the chiral SU(3) model
for the evaluation of in-medium properties of scalar, vector,
and axial-vector mesons. To find the mass modification of
the above discussed heavy mesons we will use the two-point
 correlation function $\Pi_{\mu\nu}(q)$,

$$ \Pi_{\mu\nu}(q) = i \int d^4x \ e^{iqx} \left\langle T \left\{ J_\mu(x) J^*_\nu(0) \right\} \right\rangle_{\rho\mu\nu}. \tag{2} $$

In the above equation $J_\mu(x)$ denotes the isospin averaged
current, $x = x' = (x^0, x)$ is the four coordinate, $q = q' =
(q^0, q)$ is four momentum, and $T$ denotes the time ordered
operation on the product of quantities in the brackets. From
above definition it is clear that the two-point correlation
function is actually a Fourier transform of the expectation
value of the time ordered product of two currents. The two-point
correlation function for the scalar mesons is defined as

$$ \Pi(q) = i \int d^4x \ e^{iqx} \left\langle T \left\{ J(x) J^*(0) \right\} \right\rangle_{\rho\mu\nu}. \tag{3} $$

For the scalar, vector, and axial-vector mesons isospin
average currents are given by the expressions

$$ J(x) = J^*(x) = \frac{\bar{c}(x) q(x) + \bar{q}(x) c(x)}{2}, $$

$$ J_\mu(x) = J^*_\mu(x) = \frac{\bar{c}(x) \gamma_\mu q(x) + \bar{q}(x) \gamma_\mu c(x)}{2}, \tag{4} $$

$$ J_{5\mu}(x) = J^*_{5\mu}(x) = \frac{\bar{c}(x) \gamma_5 \gamma_\mu q(x) + \bar{q}(x) \gamma_5 \gamma_\mu c(x)}{2}, $$

respectively. Note that in above equations $q$ denotes the light
$u$ or $d$ quark whereas $c$ denotes the heavy charm quark. Note
that in the present work, instead of considering the mass
splitting between particles and antiparticles, we emphasize on
the mass shift of iso-doublet $D$ and $B$ mesons as a whole and,
therefore, we consider the average in the definitions of scalar,
vector, and isovector currents which is referred as centroid
[19]. To find the mass splitting of particles and antiparticles
in the nuclear medium one has to consider the even and
odd part of QCD sum rules [18]. For example, in [18] the
mass splitting between pseudoscalar $D$ and $\bar{D}$ mesons
was investigated using the even and odd QCD sum rules whereas
in [19, 23, 24] the mass shift of $D$ mesons was investigated
under centroid approximation.

The two-point correlation function can be decomposed
into the vacuum part, a static one-nucleon part, and pion bath
contribution; that is, we can write

$$ \Pi_{\mu\nu}(q) = \Pi^0_{\mu\nu}(q) + \frac{\rho_B}{2M_N} \Pi^N_{\mu\nu}(q) + \Pi^{P.B.}_{\mu\nu}(q), \tag{5} $$

where

$$ T^N_{\mu\nu}(\omega, q) = i \int d^4x e^{iqx} \left\langle N(p) \left| T \left\{ J_\mu(x) J^*_\nu(0) \right\} \right| N(p) \right\rangle. \tag{6} $$

In the above equation $\left\langle N(p) \right|$ denotes the isospin and spin
averaged static nucleon state with the four-momentum $p =
(M_N, 0)$. The state is normalized as $\left\langle N(p') \left| N(p) \right\rangle = (2\pi)^3 2\rho_0 \delta(p - p').$ The third term, $\Pi^{P.B.}_{\mu\nu}(q)$ in (5) gives the
contribution from pion bath at finite temperature. Note that
in the present work instead of considering the contribution of
pion bath the effects of finite temperature of the nuclear
matter on the properties of $D$ and $B$ mesons will be evaluated
through the temperature dependence of scalar fields $\sigma$, $\xi$,
and $\chi$. The temperature dependence of scalar fields $\sigma$, $\xi$,
and $\chi$ modify the nucleon properties in the medium and these
modified nucleons further modify the in-medium properties
of $D$ and $B$ mesons at finite temperature and density. In
literature the properties of kaons and antikaons, $D$ mesons
and charmonium had been studied at finite temperature of
the nuclear matter using the above mentioned scalar fields $\sigma$,
$\xi$, and $\chi$ [7, 8, 12, 13, 34].

As discussed in [24], in the limit of the 3-vector $q \to 0$,
the correlation functions $T_{\mu\nu}(\omega, q)$ can be related to the $D^*N$
and $D_1N$ scattering $T$-matrices. Thus we write [24]

$$ \mathcal{S}_{D^*N}(M_{D^*}, 0) = 8\pi (M_N + M_{D^*}) a_{D^*}, \tag{7} $$

$$ \mathcal{S}_{D_1N}(M_{D_1}, 0) = 8\pi (M_N + M_{D_1}) a_{D_1}. $$

In above equation $a_{D^*}$ and $a_{D_1}$ are the scattering lengths of
$D^*N$ and $D_1N$, respectively. Similarly, we can also write the
scattering $T$ matrix corresponding to $D_0N$ ($D_0$ is a scalar
meson) in terms of the scattering lengths [23],

$$ \mathcal{S}_{D_0N}(M_{D_0}, 0) = 8\pi (M_N + M_{D_0}) a_{D_0}. \tag{8} $$

Near the pole positions of the scalar, vector, and axial-
vector mesons the phenomenological spectral densities can
be parameterized with three unknown parameters $a$, $b$, and $c$; that is, we write [19, 23, 24]

$$
\rho(\omega, 0) = \frac{\mathcal{F}_{D_1/D^*/D_1}(\omega, 0)}{\pi} \left[ \left( \omega^2 - M_{D_1/D^*/D_1}^2 + i\varepsilon \right)^2 \right]
$$

$$
+ \cdots = a \left( \omega^2 - M^2_{D_1/D^*/D_1} \right) + b \delta \left( \omega^2 - M^2_{D_1/D^*/D_1} \right) + c \theta \left( \omega^2 - s_0 \right).
$$

The term denoted by $\cdots$ represents the continuum contributions. The first term denotes the double-pole term and corresponds to the on-shell effects of the $T$-matrices,

$$
a = -8\pi \left( M_N + M_{D_1/D^*/D_1} \right) a_{D_1/D^*/D_1} f_{D_1/D^*/D_1}^2 M^2_{D_1/D^*/D_1}.
$$

(10)

Now we will write the relation between the scattering length of mesons and their in-medium mass-shift. For this we first note that the shift of squared mass of mesons can be written in terms of the parameter $a$ appearing in (9) through relation [17],

$$
\Delta m^2_{D_1/D^*/D_1} = \rho_N \frac{a}{2M_N} \left( M_N + M_{D_1/D^*/D_1} \right) a_{D_1/D^*/D_1},
$$

where in the last term we used (10). The mass shift is now defined by the relation

$$
\delta M^2_{D_1/D^*/D_1} = \sqrt{\Delta m^2_{D_1/D^*/D_1}} - m^2_{D_1/D^*/D_1}.
$$

(12)

The second term in (9) denotes the single-pole term and corresponds to the off-shell (i.e., $\omega^2 \neq M^2_{D_1/D^*/D_1}$) effects of the $T$-matrices. The third term denotes the continuum term or the remaining effects, where $s_0$ is the continuum threshold. The continuum threshold parameter $s_0$ define the scale below which the continuum contribution vanishes [42].

It can be observed from (11) and (12) that if we want to find the value of mass shift of mesons then we first need to find the value of unknown parameter $a$. For this we proceed as follows: we note that in the low energy limit, $\omega \rightarrow 0$, the $T_N(\omega, 0)$ is equivalent to the Born term $T^\text{Born}_{D_1/D^*/D_1,N}(\omega, 0)$.

We take into account the Born term at the phenomenological side,

$$
T_N(\omega^2) = T^\text{Born}_{D_1/D^*/D_1,N}(\omega^2) + \frac{a}{M^2_{D_1/D^*/D_1} - \omega^2} + \frac{b}{M^2_{D_1/D^*/D_1} - \omega^2} + \frac{c}{s_0 - \omega^2}.
$$

(13)

with the constraint

$$
\frac{a}{M^4_{D_1/D^*/D_1}} + \frac{b}{M^2_{D_1/D^*/D_1}} + \frac{c}{s_0} = 0.
$$

(14)

Note that in (13) the phenomenological side of scattering amplitude for $g_{H,N} \neq 0$ is not exactly equal to Born term but there are contributions from other terms. However, for $\omega \rightarrow 0$, $T_N$ on left should be equal to $T^\text{Born}_N$ on right side of (13) and this requirement results in constraint given in (14). As we will discuss below the constraint (14) help in eliminating the parameter $c$ and scattering amplitude will be function of parameters $a$ and $b$ only. The Born terms to be used in (13) for scalar, vector, and axial-vector mesons are given by following relations [23, 24]:

$$
T^\text{Born}_{D_1,N}(\omega, 0) = \frac{2f_{D_1}^2 M^2_{D_1} M_N (M_H - M_N) g_{D_1,N}^2}{\left( \omega^2 - (M_H - M_N)^2 \right) \left( \omega^2 - M^2_{D_1} \right)}.
$$

$$
T^\text{Born}_{D_N}(\omega, 0) = \frac{2f_{D_N}^2 M^2_{D_N} M_N (M_H + M_N) g_{D_N,N}^2}{\left( \omega^2 - (M_H + M_N)^2 \right) \left( \omega^2 - M^2_{D_N} \right)}.
$$

(15)

In the above equations $g_{D_1,N}$, $g_{D_N,N}$, and $g_{N,N}$ are the coupling constants. $M_H$ is the mass of the hadron; for example, corresponding to charm mesons we have $\Lambda_c$ and $\Sigma_c$, whereas corresponding to bottom mesons we have the hadrons $\Lambda_b$ and $\Sigma_b$. Corresponding to charm mesons we take the average value of the masses of $M_{\Lambda_c}$ and $M_{\Sigma_c}$ and it is equal to 2.4 GeV. For the case of mesons having bottom quark, $b$, we consider the average value of masses of $\Lambda_b$ and $\Sigma_b$ and it is equal to 5.7 GeV.

Now we write the equation for the Borel transformation of the scattering matrix on the phenomenological side and equate that to the Borel transformation of the scattering matrix for the operator expansion side. For the scalar meson, $D_0$, the Borel transformation equation is written as [23]

$$
a \left\{ \frac{1}{M^2} e^{-M^2_{b} / M^2} - \frac{s_0}{M^2_{b}} e^{-s_0 / M^2} \right\}
$$

$$
+ b \left\{ e^{-M^2_{b} / M^2} - \frac{s_0}{M^2_{b}} e^{-s_0 / M^2} \right\}
$$

$$
+ A \left[ \frac{1}{(M_H - M_N)^2 - M^2_{b}} - \frac{1}{M^2} \right] e^{-M^2_{b} / M^2}
$$

$$
- \frac{A}{(M_H - M_N)^2 - M^2_{b}} e^{-(M_H - M_N)^2 / M^2}
$$

$$
= \left\{ \frac{m^2 \langle \bar{q}q \rangle_N}{2} - \langle \bar{q}^4 iD_0 q \rangle_N + \frac{m^2 \langle \bar{q}^4 iD_0 q \rangle_N}{M^2} \right\} e^{-m^2_{b} / M^2}
$$

$$
+ \frac{m \langle q^4 \rangle_N}{M^2} + \frac{m^2 \langle \bar{q}^4 iD_0 q \rangle_N}{M^2} e^{-m^2_{b} / M^2}
$$

$$
+ \frac{m^4 \langle \bar{q}^2 q \rangle_N}{2M^2} + \frac{m^2 \langle \bar{q}^4 iD_0 q \rangle_N}{M^2} e^{-m^2_{b} / M^2}
$$

$$
+ \frac{m^6 \langle q^2 iD_0 q \rangle_N}{M^2} e^{-m^2_{b} / M^2}.
$$

$$
= \left\{ \frac{m^2 \langle \bar{q}q \rangle_N}{2} - \langle \bar{q}^4 iD_0 q \rangle_N + \frac{m^2 \langle \bar{q}^4 iD_0 q \rangle_N}{M^2} \right\} e^{-m^2_{b} / M^2}
$$

$$
+ \frac{m \langle q^4 \rangle_N}{M^2} + \frac{m^2 \langle \bar{q}^4 iD_0 q \rangle_N}{M^2} e^{-m^2_{b} / M^2}
$$

$$
+ \frac{m^4 \langle \bar{q}^2 q \rangle_N}{2M^2} + \frac{m^2 \langle \bar{q}^4 iD_0 q \rangle_N}{M^2} e^{-m^2_{b} / M^2}
$$

$$
+ \frac{m^6 \langle q^2 iD_0 q \rangle_N}{M^2} e^{-m^2_{b} / M^2}.
$$
where \( A = 2f_{D_0}^2 M_{D_0}^2 N_0 (M_H - M_N) g_{D_0 NH}^2 ((M_H - M_N)^2 - M_{D_0}^2) \).

Note that in (16) we have two unknown parameters \( a \) and \( b \). We differentiate (16) w.r.t. \( 1/M^2 \) so that we could have two equations and two unknowns. By solving those two coupled equations we will be able to get the values of parameters \( a \) and \( b \). Same procedure will be applied to obtain the values of parameters \( a \) and \( b \) corresponding to vector and axial-vector mesons. For vector meson, \( D^* \), the Borel transformation equation is given by [24]

\[
\frac{1}{2} \left\{ \frac{1}{M^2} e^{-M_{D^*}^2/M^2} - \frac{s_0}{M_{D^*}^2} e^{-s_0/M^2} \right\}
+ b \left\{ e^{-M_{D^*}^2/M^2} - \frac{s_0}{M_{D^*}^2} e^{-s_0/M^2} \right\} + B \left[ \frac{1}{(M_H + M_N)^2 - M_{D^*}^2} - \frac{1}{M^2} \right] e^{-M_{D^*}^2/M^2}
- \frac{B}{(M_H + M_N)^2 - M_{D^*}^2} e^{-(M_H + M_N)^2/M^2}
= \left\{ -\frac{m_c \langle \bar{q} q \rangle_N}{2} - \frac{2 \langle q^1 i D_0 q \rangle_N}{3} + \frac{m_c^2 \langle q^1 i D_0 q \rangle_N}{M^2} \right\} e^{-M_{D^*}^2/M^2}
\times e^{-m_\chi^2/M^2} + \frac{m_c \langle \bar{q} q, \sigma G q \rangle_N}{3 M^2} e^{-m_\chi^2/M^2}
+ \left\{ \frac{8 m_c \langle \bar{q} i D_0 i D_0 q \rangle_N}{3 M^2} - \frac{m_c^3 \langle q i D_0 i D_0 q \rangle_N}{M^4} \right\} e^{-m_\chi^2/M^2}
- \frac{1}{24} \left\{ \frac{\alpha_s G G}{\pi} \right\} N_0 \int_0^1 dx \left( 1 + \frac{m_c^2}{2 M^2} \right) e^{-m_\chi^2/M^2}
+ \frac{1}{48 M^2} \left\{ \frac{\alpha_s G G}{\pi} \right\} N_0 \int_0^1 dx \left( 1 - \frac{x}{x} \left( \frac{m_c^2}{M^2} - \frac{m_\chi^2}{M^2} \right) \right) e^{-m_\chi^2/M^2},
\]

where \( B = 2f_{D_0}^2 M_{D_0}^2 N_0 (M_H + M_N) g_{D_0 NH}^2 ((M_H + M_N)^2 - M_{D_0}^2) \). For the axial-vector meson, \( D_1 \), the Borel transformation equation is given by [24]

\[
\frac{1}{M^2} e^{-M_{D_1}^2/M^2} - \frac{s_0}{M_{D_1}^2} e^{-s_0/M^2} \}
+ b \left\{ e^{-M_{D_1}^2/M^2} - \frac{s_0}{M_{D_1}^2} e^{-s_0/M^2} \right\}

+ C \left[ \frac{1}{(M_H - M_N)^2 - M_{D_1}^2} - \frac{1}{M^2} \right] e^{-M_{D_1}^2/M^2}
- \frac{C}{(M_H - M_N)^2 - M_{D_1}^2} e^{-(M_H - M_N)^2/M^2}
= \left\{ \frac{m_c \langle \bar{q} q \rangle_N}{2} - \frac{2 \langle q^1 i D_0 q \rangle_N}{3} + \frac{m_c^2 \langle q^1 i D_0 q \rangle_N}{M^2} \right\} e^{-m_\chi^2/M^2}
- \frac{m_c \langle \bar{q} q, \sigma G q \rangle_N}{3 M^2} e^{-m_\chi^2/M^2}
+ \left\{ \frac{8 m_c \langle \bar{q} i D_0 i D_0 q \rangle_N}{3 M^2} - \frac{m_c^3 \langle q i D_0 i D_0 q \rangle_N}{M^4} \right\} e^{-m_\chi^2/M^2}
- \frac{1}{24} \left\{ \frac{\alpha_s G G}{\pi} \right\} N_0 \int_0^1 dx \left( 1 + \frac{m_c^2}{2 M^2} \right) e^{-(M_H - M_N)^2/M^2}
+ \frac{1}{48 M^2} \left\{ \frac{\alpha_s G G}{\pi} \right\} N_0 \int_0^1 dx \left( 1 - \frac{x}{x} \left( \frac{m_c^2}{M^2} - \frac{m_\chi^2}{M^2} \right) \right) e^{-(M_H - M_N)^2/M^2},
\]

where \( C = 2f_{D_0}^2 M_{D_0}^2 N_0 (M_H - M_N) g_{D_0 NH}^2 ((M_H - M_N)^2 - M_{D_0}^2) \). In the above equations \( m_\chi^2 = m_c^2 / x \).

As discussed earlier, in determining the properties of hadrons from QCD sum rules, we will use the values of quark and gluon condensates as calculated using chiral SU(3) model. Any operator \( \mathcal{O} \) on OPE side can be written as [17, 42, 43]

\[
\mathcal{O}_{p\chi} = \mathcal{O}_{\text{vacuum}} + 4 \int \frac{d^3p}{(2\pi)^3} 2E_p \langle 0 | \mathcal{O} | n (p) \rangle \langle n (p) | \mathcal{O} | 0 \rangle
+ 3 \int \frac{d^3k}{(2\pi)^3} 2E_k \langle 0 | \mathcal{O} | n_b (k) \rangle \langle n_b (k) | \mathcal{O} | 0 \rangle
\]

\[
= \mathcal{O}_{\text{vacuum}} + \frac{P_B}{2 M} \mathcal{O}_{N} + \mathcal{O}_{P.B.}
\]

In the above equation, \( \mathcal{O}_{p\chi} \) gives us the expectation value of the operator at finite baryonic density. The term \( \mathcal{O}_{\text{vacuum}} \) stands for the vacuum expectation value of the operator, \( \mathcal{O}_N \) gives us the nucleon expectation value of the operator, and \( \mathcal{O}_{P.B.} \) denotes the contribution from the pion bath at finite temperature. Consider \( n_B = [e^{E_B/T} - 1]^{-1} \) and \( n_F = [e^{E_F/T} - 1]^{-1} \) are the thermal Boson and Fermion distribution functions. Within the chiral SU(3) model the quark and gluon condensates can be expressed in terms of scalar fields \( \sigma \), \( \zeta \), and \( \chi \). As discussed earlier, the finite temperature effects in the present investigation will be evaluated through the scalar fields and, therefore, contribution of third term will not be considered. However, for completeness we will compare the temperature dependence of scalar quark and scalar gluon condensates at zero baryon density as evaluated in the present work with the situation when the temperature dependence is evaluated using only pion bath contribution.
Thus within chiral SU(3) model, we can find the values of $\bar{\sigma}_N$ at finite density of the nuclear medium and hence can find $\bar{\sigma}_N$ using

$$\bar{\sigma}_N = \left[ \bar{\sigma}_N - \bar{\sigma}_{\text{vacuum}} \right] \frac{2M_N}{\rho_0}. \tag{20}$$

The quark condensate, $\langle \bar{q}q \rangle$, can be extracted from the explicit symmetry breaking term of the Lagrangian density and is given by

$$\sum_i m_i \langle \bar{q}_i q_i \rangle = -\mathcal{L}_{SB} = \left( \frac{\chi}{\chi_0} \right)^2 \left[ m_{\pi}^2 f_{\pi} \sigma + \left( \sqrt{2} m_{\pi}^2 f_{\pi} - \frac{1}{\sqrt{2}} m_{\pi}^2 f_{\pi} \right) \zeta \right]. \tag{21}$$

In our present investigation of hadron properties, we are interested in light quark condensates, $\bar{u}u$ and $\bar{d}d$, which are proportional to the nonstrange scalar field $\sigma$ within chiral SU(3) model. Considering equal mass of light quarks, $u$ and $d$ that is, $m_u = m_d = m_q = 0.006 \text{ GeV}$, we can write,

$$\langle \bar{q}q \rangle_{\rho_0} = \frac{1}{2m_q} \left( \frac{\chi}{\chi_0} \right)^2 \left[ m_{\sigma}^2 f_{\pi} \sigma \right]. \tag{22}$$

The condensate $\langle \bar{q}g_\sigma Gq \rangle_{\rho_0}$ is given by the following [44]:

$$\langle \bar{q}g_\sigma Gq \rangle_{\rho_0} = \lambda^2 \langle \bar{q}q \rangle_{\rho_0} + 3.0 \text{ GeV}^2 \rho_0. \tag{23}$$

Also we write [44]

$$\langle \bar{q}D_0iD_0q \rangle_{\rho_0} + \frac{1}{8} \langle \bar{q}g_\sigma Gq \rangle_{\rho_0} = 0.3 \text{ GeV}^2 \rho_0. \tag{24}$$

As discussed above the quark condensate, $\langle \bar{q}q \rangle_{\rho_0}$, can be calculated within the chiral SU(3) model. This value of $\langle \bar{q}q \rangle_{\rho_0}$ can be used through (23) and (24) to calculate the value of condensates $\langle \bar{q}g_\sigma Gq \rangle_{\rho_0}$ and $\langle \bar{q}D_0iD_0q \rangle_{\rho_0}$ within chiral SU(3) model. The value of condensate $\langle q^i D_i q \rangle$ is equal to 0.18 GeV$^2 \rho_0$ [44].

At finite temperature and zero baryon density we can write the expectation values of quark condensates and scalar gluon condensates as [42]

$$\langle \bar{q}q \rangle_T = \langle \bar{q}q \rangle_{\text{vacuum}} \left[ 1 - \frac{T^2}{8f_{\pi}^2} B_1 \left( \frac{m_\pi}{T} \right) \right], \tag{25}$$

$$\langle \bar{q}g_\sigma G_{\mu\nu}G^{\mu\nu} \rangle_T = \frac{\alpha_s}{\pi} \langle \bar{q}g_\sigma G_{\mu\nu}G^{\mu\nu} \rangle_{\text{vacuum}} \left[ 1 - \frac{1}{9} m_{\pi}^2 T^2 B_1 \left( \frac{m_\pi}{T} \right) \right], \tag{26}$$

respectively. In (25) and (26), $B_1(x) = (6/n^2) \int_0^\infty dy y^{3/2} - x^2/(e^y - 1)$. From (25) and (26) we observe that the contribution from pion bath to the expectation values of operator arises only at finite temperature. Note that in (16), (17), and (18) we need the nucleon expectation values of various condensates which can be evaluated in general using (20). In (1), (22), (23), and (24) the values of condensates are given at finite value of baryonic density. To find the corresponding nucleon expectation values of various condensates we use the values of condensates at finite baryonic density from (1), (22), (23), and (24) in (20). The equation (20) is then further used in (16), (17), and (18) for calculation of medium modifications of $D$ mesons.

It may be noted that the QCD sum rules for the evaluation of in-medium properties of scalar mesons, $B_0$, vector mesons, $B^*$, and axial-vector mesons, $B_1$, can be written by replacing masses of charmed mesons, $D_{bc}, D^*, \text{ and } D_1$, by corresponding masses of bottom mesons $B_0, B^*, \text{ and } B_1$ in (16), (17), and (18), respectively. Also the bare charm quark mass, $m_c$, will be replaced by the mass of bottom quark, $m_b$.

4. Results and Discussions

In this section we will present the results of our investigation of in-medium properties of scalar ($D_0, B_0$), vector ($D^*, B^*$), and axial-vector mesons ($D_1, B_1$). The nuclear matter saturation density used in the present investigation is 0.15 fm$^{-3}$. The values of various coupling constants $g_{D_{bc}N\Lambda} \approx g_{D_{bc}N\Sigma} \approx g_{B_{bc}N\Lambda} \approx g_{B_{bc}N\Sigma}$ are approximated to 6.74 [23]. The coupling constants $g_{D^*N\Lambda} \approx g_{D^*N\Sigma} \approx g_{D_{bc}N\Lambda} \approx g_{D_{bc}N\Sigma} \approx g_{B_{bc}N\Lambda} \approx g_{B_{bc}N\Sigma}$ are approximated to 3.86 [24]. The masses of mesons $M_{D_{bc}}, M_{B_{bc}}, M_{D^*}, M_{D_{bc}}, M_{D_{bc}},$ and $M_{B_{bc}}$ to be used in present investigation are 2.355, 5.74, 2.01, 5.325, 2.42, and 7.57 GeV, respectively. The values of decay constants $f_{D_{bc}}, f_{B_{bc}}, f_{D^*}, f_{B^*},$ and $f_{B_{bc}}$ are 0.334, 0.28, 0.270, 0.195, 0.305, and 0.255 GeV, respectively. The values of threshold parameters, $s_0$, corresponding to $D_0, B_0, D^*, B^*, D_1, \text{ and } B_1$ mesons are 8, 39, 6.5, 35, 8.5, and 39 GeV$^2$, respectively [23, 24]. As discussed earlier the mass-shift of scalar, vector, and axial-vector mesons is calculated through the parameter $\alpha$ which is related to scattering length through (10) to (12). This parameter $\alpha$, for example, for $D^*$, is calculated by solving the coupled equations as discussed after (16) and is subjected to the medium modifications through the medium dependence of condensates. The medium dependence of condensates is further evaluated through the scalar fields $\sigma, \zeta$, and $\chi$. The various coupling constants listed above, the decay constants of $D$ and $B$ mesons, and the threshold parameter $s_0$ are not subjected to the medium modifications. In the present work we will show the variation of mass shift as a function of squared Borel mass parameter, $M$. The Borel window is chosen such that there is almost no change in the mass of $D$ and $B$ mesons w.r.t variation in Borel mass parameter. For the charmed scalar, $D_{bc}, \text{ and } D^*$, and axial-vector, $D_1$ mesons Borel windows are found to be (6.1–7.4), (4.5–5.4), and (6.5–7.6) GeV$^2$, respectively. The Borel windows for bottom scalar, $B_{bc}, \text{ vector, } B^*, \text{ and } B_1$ mesons are (33–39), (22–24), and (34–37) GeV$^2$, respectively.

In the present work we are studying the in-medium masses of scalar, vector, and axial-vector mesons using QCD sum rules and chiral SU(3) model. Since we are evaluating the
quark and gluon condensates within the chiral SU(3) model through the modification of scalar fields $\sigma$ and $\zeta$ and the scalar dilaton field $\chi$, so first we will discuss in short the effect of temperature and density of the medium on the values of scalar fields $\sigma$ and $\zeta$ and the scalar dilaton field $\chi$. We observe that as a function of density of the nuclear medium the magnitude of scalar fields decreases. For example, at nuclear saturation density, $\rho_B = \rho_0$, the drop in magnitude of the scalar fields $\sigma$ and $\zeta$ and the dilaton field $\chi$ is observed to be 3.49, 9.54 and 3.39 MeV, respectively. At baryon densities, $4\rho_0$, these values change to 63.41, 14.51 and 13.06 MeV, respectively.

At zero baryon density, we observe that the magnitude of the scalar fields $\sigma$ and $\zeta$ and the dilaton field $\chi$ decreases with the increase in the temperature. However, the change in the values of scalar fields with temperature of the medium is observed to be very small. The reason for the nonzero values of scalar fields at finite temperature and zero baryon density of the medium is the formation of baryon-antibaryon pairs [36, 45, 46]. At finite baryon densities, the magnitude of the scalar fields increases with increase in the temperature of symmetric nuclear medium. This also leads to the increase in the masses of the nucleons with the temperature of the nuclear medium for finite baryon densities [12]. At nuclear saturation density, $\rho_B = \rho_0$, the magnitude of the scalar fields, $\sigma$ and $\zeta$, and the dilaton field $\chi$ increases by 5.47, 1.28 and 0.96 MeV, respectively, as we move from $T = 0$ to $T = 150$ MeV, respectively.

In Figures 1 and 2 we show the variation of the light scalar quark condensate $\bar{q}q$, given by (22) and the scalar gluon condensate $G_0 = \langle (\alpha_s/\pi) G^a_\mu G^{a\mu} \rangle$, given by (1), respectively, as a function of density of the symmetric nuclear medium. We show the results for temperatures, $T = 0, 50, 100, \text{and } 150$ MeV, respectively. From (22) we observe that the value of the scalar quark condensate $\bar{q}q$ is directly proportional to the scalar-isoscalar field, $\sigma$. Therefore, the behavior of the $\bar{q}q$ as a function of temperature and density of the nuclear medium will be the same as that of $\sigma$ field. For given value of temperature of the nuclear medium the magnitude of the light quark condensate decreases with the increase in the density. For example, at $T = 0$, the values of light quark condensate are observed to be $-0.8829 \times 10^{-2}$ and $-0.4200 \times 10^{-2}$ GeV$^3$ at baryon densities, $\rho_B = \rho_0$ and $4\rho_0$, respectively. At temperature, $T = 150$ MeV, these values of quark condensates changes to $-0.9681 \times 10^{-2}$ and $-0.4993 \times 10^{-2}$ GeV$^3$, respectively. At zero baryon density the magnitude of the $\bar{q}q$ decreases with the increase in the temperature of the nuclear medium. At $\rho_B = 0$, the values of $\bar{q}q$ are observed to be $-1.4014 \times 10^{-2}$, $-1.4014 \times 10^{-2}$, $-1.4006 \times 10^{-2}$, and $-1.3628 \times 10^{-2}$ GeV$^3$ at temperatures $T = 0, 50, 100, \text{and } 150$ MeV, respectively.

From Figure 2, we observe that the values of the scalar gluon condensates decrease with the increase in the density of the nuclear medium. At baryon density, $\rho_B = \rho_0$, the values of $G_0$ are observed to be $1.90646 \times 10^{-2}$ GeV$^4$, $1.9119 \times 10^{-2}$ GeV$^4$, $1.91755 \times 10^{-2}$ GeV$^4$, and $1.92 \times 10^{-2}$ GeV$^4$ for temperatures, $T = 0, 50, 100, \text{and } 150$ MeV, respectively. For the same values of the temperature, in the absence of finite quark masses, the values of $G_0$ are observed to be $2.269 \times 10^{-2}$ GeV$^4$, $2.2771 \times 10^{-2}$ GeV$^4$, $2.2857 \times 10^{-2}$ GeV$^4$, and $2.29 \times 10^{-2}$ GeV$^4$ for $\rho_B = \rho_0$. For baryon density, $\rho_B = 4\rho_0$, the values of $G_0$ are given as $1.7367 \times 10^{-2}$ GeV$^4$ ($2.06 \times 10^{-2}$ GeV$^4$), $1.74612 \times 10^{-2}$ GeV$^4$ ($2.0713 \times 10^{-2}$ GeV$^4$), $1.7656 \times 10^{-2}$ GeV$^4$ ($2.094 \times 10^{-2}$ GeV$^4$), and $1.78 \times 10^{-2}$ GeV$^4$ ($2.112 \times 10^{-2}$ GeV$^4$) for values of temperature, $T = 0, 50, 100, \text{and } 150$ MeV, respectively, for the cases of the finite (zero) quark masses in the trace anomaly.

It may be noted that in the above discussion the finite temperature effects of the nuclear medium on the values of quark and gluon condensates are evaluated through the temperature dependence of scalar fields $\sigma$ and $\zeta$ and the dilaton field $\chi$. 

![Figure 1: (Color online) The light quark condensate $\bar{q}q$ plotted as a function of density of the nuclear medium, in units of nuclear saturation density, for different values of temperatures ($T = 0, 50, 100, \text{and } 150$ MeV).](image1)

![Figure 2: (Color online) The scalar gluon condensate $G_0$ plotted as a function of density of the nuclear medium, in units of nuclear saturation density, for different values of temperatures ($T = 0, 50, 100, \text{and } 150$ MeV).](image2)
In literature the scalar quark and gluon condensates at finite temperature are evaluated due to contribution from pion bath using (25) and (26), respectively [42]. Using $m_0(q^2) = -0.11 \text{ GeV}^2$ and $(\alpha_s/r_G)^2 = 0.005 \text{ GeV}^4$ [42] we calculate the quark and gluon condensates at finite temperature and zero baryon density using (25) and (26), respectively. The values of scalar quark condensates are observed to be $-1.824 \times 10^{-2}$, $-1.751 \times 10^{-2}$, and $-1.566 \times 10^{-2} \text{ GeV}^3$ at temperatures $T = 50, 100,$ and $150 \text{ MeV}$, respectively. These values of scalar quark condensates can be compared to the values $1.4014 \times 10^{-2}$, $-1.4006 \times 10^{-2}$, and $-1.3628 \times 10^{-2} \text{ GeV}^3$ at $T = 50, 100,$ and $150 \text{ MeV}$, respectively. However, for zero baryon density, the values of scalar quark condensates are found to vary effectively with temperature above critical temperature $T_C = 145.8 \text{ MeV}$ at baryonic density $\rho_B = \rho_0$ and zero temperature, the values of scalar quark and gluon condensates are observed to be $-0.6222 \times 10^{-2} \text{ GeV}^3$ and $1.9869 \times 10^{-2} \text{ GeV}^4$, respectively. These values of scalar quark and gluon condensates are observed to be $-0.8829 \times 10^{-2} \text{ GeV}^3$ and $1.9065 \times 10^{-2} \text{ GeV}^4$, respectively which were calculated without the medium modification of $f_\sigma, f_K, m_\pi,$ and $m_K$. We conclude that the medium modification of $f_\sigma, f_K, m_\pi,$ and $m_K$ causes more decrease in the values of scalar quark and gluon condensates at finite baryonic density.

Now we will calculate the in-medium masses of scalar, vector, and axial-vector mesons using the values of condensates from chiral SU(3) model. In Figure 3, the subfigures (a), (c), and (e) show the variation of the mass shift of scalar mesons $D_0$ as a function of square of the Borel mass parameter, $M$, at nuclear matter densities $\rho_0, 2\rho_0,$ and $4\rho_0$, respectively. The subfigures (b), (d), and (f) show the variation of the scalar mesons $B_0$ as a function of square of the Borel mass parameter, $M$, at nuclear matter densities $\rho_0, 2\rho_0,$ and $4\rho_0$, respectively. In each subplot we have shown the results at temperatures $T = 0, 50, 100,$ and $150 \text{ MeV}$. We observe that for scalar mesons, $D_0$, at temperature, $T = 0$, the values of mass shift are found to be $76, 114,$ and $148 \text{ MeV}$ at baryon densities $\rho_0, 2\rho_0,$ and $4\rho_0$, respectively. For temperature $T = 50 \text{ MeV}$ the values of mass shift are observed to be $71, 109,$ and $144 \text{ MeV}$ at baryon densities $\rho_0, 2\rho_0,$ and $4\rho_0$, respectively. At temperature $T = 100 \text{ MeV}$ the above values of mass shift are observed to be $66, 103,$ and $139 \text{ MeV}$, whereas at $T = 150 \text{ MeV}$ the values of mass shift changes to $58, 94,$ and $131 \text{ MeV}$ at baryon densities $\rho_0, 2\rho_0,$ and $4\rho_0$, respectively. From the above discussion we conclude that for a given value of temperature the mass shift of scalar mesons, $D_0$, increases as a function of density of the nuclear medium. On the other hand as a function of temperature of the nuclear medium, for a constant value of density, the mass shift of scalar mesons, $D_0$, decreases.

As we can see from Figure 3 the values of mass shift of scalar $B_0$ mesons also increases as a function of density of the nuclear medium. At temperature, $T = 0$, the values of mass shift are observed to be $224, 334,$ and $420 \text{ MeV}$ at baryon densities $\rho_0, 2\rho_0,$ and $4\rho_0$, respectively. At temperature $T = 50 \text{ MeV}$ the above values of mass shift are found to be $211, 321,
Figure 3: (Color online) Figure shows the variation of the mass shift of scalar mesons $D_0$ (subplots (a), (c), and (e)) and $B_0$ (subplots (b), (d), and (f)) as a function of the squared Borel mass parameter, $M^2$. We show the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. For each value of baryon density the results are shown at temperatures, $T = 0, 50, 100, \text{ and } 150 \text{ MeV}$. Also the medium modifications of decay constants and masses of pions and kaons are considered while evaluating the above shown mass-shift of scalar mesons.
and 413 MeV at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$, respectively. For temperature $T = 100$ MeV, the values of mass shift are found to be 196, 302, and 399 MeV whereas at $T = 150$ MeV, these values of mass shift changes to 171, 274, and 372 MeV at $\rho_0$, $2\rho_0$, and $4\rho_0$, respectively. For a given value of nuclear matter density, the values of mass shift present for the $B_0$ mesons are found to decrease with the increase in the temperature of the medium.

We also calculate the values of scattering lengths of scalar mesons using (10) for different values of density and temperature of the medium. At temperature $T = 0$ and baryon densities, $p_B = \rho_0$ and $4\rho_0$, the values of scattering lengths for scalar mesons $D_0$ ($B_0$) are observed to be 1.42 (5.05) and 0.70 (2.41) fm, respectively. For temperature $T = 100$ MeV and baryon densities, $p_B = \rho_0$ and $4\rho_0$, the values of scattering lengths for $D_0$ ($B_0$) are observed to be 1.23 (4.40) and 0.66 (2.28) fm, respectively. We note that the value of scattering length decreases as a function of density and temperature of the nuclear medium. As discussed above the scattering lengths are evaluated using (10). In this equation we have the parameter $a$ which is directly proportional to the scattering length of mesons. As discussed earlier the value of parameter $a$ is evaluated by solving simultaneously (16) and the equation obtained by differentiate (16) w.r.t. $1/M^2$. The magnitude of parameter $a$ is found to decrease as we move from low to higher value of baryonic density or from zero to finite value of temperature of the medium. This behavior of parameter $a$ as a function of density and temperature of the medium results in the similar changes in the values of scattering lengths of scalar mesons.

Figure 4 shows the variation of the mass shift of vector mesons $D^*$ and $B^*$ as a function of square of the Borel mass parameter. Here also we have shown the results at nuclear matter densities $\rho_0$, $2\rho_0$, and $4\rho_0$. We observe that at nuclear matter density $\rho_B = \rho_0$, the values of mass shift for vector mesons, $D^*$ are observed to be $-76$, $-71$, $-65$, and $-56$ MeV at temperatures $T = 0$, $50$, $100$, and $150$ MeV, respectively. For baryon density $p_B = 2\rho_0$ the values of mass shift are found to be $-111$, $-106$, $-98$, and $-87$ MeV, whereas for $p_B = 4\rho_0$ the values of the mass shift changes to $-128$, $-125$, $-120$, and $-108$ MeV at temperatures $T = 0$, $50$, $100$, and $150$ MeV, respectively. Note that for a given value of density, the magnitude of the mass shift of vector mesons decreases as a function of temperature of the nuclear medium. On the other hand as a function of density of the medium the magnitude of the mass shift of the vector mesons $D^*$ increases. For nuclear matter saturation density the values of mass shift for the $B^*$ mesons are found to be $-366$, $-344$, $-311$, and $-275$ MeV at temperatures $T = 0$, $50$, $100$, and $150$ MeV, respectively. At baryon density $2\rho_0$ the above values of mass shift change to $-557$, $-534$, $-498$, and $-447$ MeV at temperatures $T = 0$, $50$, $100$, and $150$ MeV, respectively. For baryon density $4\rho_0$ the above values of mass shift are found to be $-701$, $-689$, $-662$, and $-609$ MeV at temperatures $T = 0$, $50$, $100$, and $150$ MeV, respectively. At temperature $T = 0$ and baryon densities, $p_B = \rho_0$ and $4\rho_0$, the values of scattering lengths for vector mesons $D^*$ ($B^*$) are observed to be $-1.31$ ($-7.72$) and $-0.54$ ($-3.58$) fm, respectively. For temperature $T = 100$ MeV and baryon densities, $p_B = \rho_0$ and $4\rho_0$, the values of scattering lengths for $D^*$ ($B^*$) are observed to be $-1.12$ ($-6.72$) and $-0.51$ ($-3.54$) fm, respectively. We observe that the magnitude of the scattering lengths of vector mesons decreases in moving from low to higher value of density or temperature of the medium.

In Figure 5, for given values of temperatures and densities we have shown the variation of mass shift of axial-vector mesons $D_1$ and $B_1$ as a function of square of the Borel mass parameter. We observe that at baryon density $p_B = \rho_0$, the values of mass shift for axial-vector meson $D_1$ are observed to be 73, 69, 63, and 55 MeV at temperatures, $T = 0$, $50$, $100$, and $150$ MeV. At baryon density $2\rho_0$ ($4\rho_0$) the values of mass shift are found to be $108$ ($131$), $104$ ($128$), $97$ ($123$), and $87$ ($113$) MeV at temperatures, $T = 0$, $50$, $100$, and $150$ MeV. For the axial-vector meson $B_1$ at baryon density $p_B = (2\rho_0)$ the values of mass shift are found to be 267 (396), 251 (381), 233 (357), and 203 (324) MeV at temperatures, $T = 0$, $50$, $100$, and $150$ MeV, respectively. At baryon density $4\rho_0$, the values of mass shift are found to be 492, 485, 467, and 434 MeV at temperatures, $T = 0$, $50$, $100$, and $150$ MeV, respectively. The values of scattering lengths for axial-vector mesons $D_1$ ($B_1$) at temperature $T = 0$ and baryon densities, $p_B = \rho_0$ and $4\rho_0$, are observed to be 1.38 (6.02) and 0.62 (2.83) fm, respectively. For temperature $T = 100$ MeV and baryon densities, $p_B = \rho_0$ and $4\rho_0$, the values of scattering lengths for $D^*$ ($B^*$) changes to 1.19 ($5.25$) and $0.58$ ($2.69$) fm, respectively. From the above discussions we observe a positive value of mass shift for scalar ($D_0$, $B_0$) and axial-vector mesons ($D_1$, $B_1$) in the nuclear medium. However, the values of mass shift for vector mesons ($D^*$, $B^*$) are found to be negative. It means the masses of above scalar and axial-vector mesons in the nuclear medium may be large compared to the value in free space and this may lead to a decrease in the yield of these mesons in heavy-ion collisions.

Now we will discuss the effect of different individual condensates on the in-medium modification of scalar ($D_0$, $B_0$), vector ($D^*$, $B^*$), and axial-vector ($D_1$, $B_1$) mesons. In Figures 6, 7, and 8 we compare the contributions of individual condensates to the mass shift of scalar mesons, $D_0$, vector mesons, $D^*$, and axial-vector mesons, $D_1$, respectively. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. Note that in Figure 6 and in the subsequent figures of this paper the word “Total” is for the contribution of all condensates to the properties of mesons. Similarly, “Quark$q_i$”, “Quark$q_2$”, “Quark$q_3$” and “Gluon$q_4$” are denoting the contribution of $(q_i q_i)$, $(q_i^1 D_0 q_i)$, $(q_i D_0^1 D_0 q_i)$, and $(a_i GG/\pi)$, respectively, to the in-medium properties of mesons.

We observe that at temperature $T = 0$; if we consider the contribution of scalar quark condensates, only then the values of mass shift for the scalar $D_0$ mesons are found to be $73.30$ and $150.68$ MeV at nuclear matter density $p_B = \rho_0$ and $4\rho_0$, respectively. When we consider the individual contributions of $(q_i^1 D_0 q_i)$, $(q_i D_0^1 D_0 q_i)$, and $(a_i GG/\pi)$ condensates then the values of mass shift at $p_B = (4\rho_0)$ and $T = 0$ are observed to be $4.01$ ($16.07$), $12.70$ ($37.87$), and $6.68$ ($25.96$) MeV,
Figure 4: (Color online) Figure shows the variation of the mass shift of vector mesons $D^*$ (subplots (a), (c), and (e)) and $B^*$ (subplots (b), (d), and (f)) as a function of the squared Borel mass parameter, $M^2$. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. For each value of the baryon density the results are shown at temperatures, $T = 0, 50, 100,$ and $150$ MeV.
Figure 5: (Color online) Figure shows the variation of the mass shift of axial-vector mesons $D_1$ (subplots (a), (c), and (e)) and $B_1$ (subplots (b), (d), and (f)) as a function of the squared Borel mass parameter, $M^2$. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. For each value of the baryon density the results are shown at temperatures, $T = 0, 50, 100, \text{and} 150 \text{MeV}$. 
Figure 6: (Color online) Figure shows the contribution of individual condensates to the mass shift of scalar mesons $D_0$ as a function of the squared Borel mass parameter, $M^2$. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. 
Figure 7: (Color online) Figure shows the contribution of individual condensates to the mass shift of vector mesons $D^*$ as a function of the squared Borel mass parameter, $M^2$. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. 

\begin{align*}
\text{Total} & \quad \text{Quark}_1 \\
\text{Gluon}_1 & \quad \text{Quark}_2 \\
\text{Quark}_4 & \quad \text{Quark}_2 = 0
\end{align*}
Figure 8: (Color online) Figure shows the contribution of individual condensates to the mass shift of axial-vector mesons $D_1$ as a function of the squared Borel mass parameter, $M^2$. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. 
respectively. From above discussion we observe that the maximum contribution to the in-medium modification of scalar $D_0$ mesons is from light quark condensate $\langle \bar{q} q \rangle$. Note that leaving the quark condensate, $\langle q^4 D_0 q^4 \rangle$, all the other condensates have been evaluated within chiral SU(3) model in the present investigation. So in Figure 6 we also show the variation of the mass shift of scalar $D_0$ mesons as a function of squared Borel mass parameter when we neglect the contribution of $\langle q^4 D_0 q^4 \rangle$ condensate. We observe that if we neglect the contribution of $\langle q^4 D_0 q^4 \rangle$ condensate then the values of mass shift are found to be 78.70 and 158.99 MeV at densities $\rho_0$ and $4\rho_0$, respectively. Note that considering the contribution of all condensates, at temperature $T = 0$, the values of mass shift were 75.88 and 148.05 MeV at densities $\rho_0$ and $4\rho_0$, respectively. So we observe that if we neglect the condensate $\langle q^4 D_0 q^4 \rangle$ then there is a percentage change of 4% and 7% in the mass shift of $D_0$ mesons at densities $\rho_0$ and $4\rho_0$, respectively. From Figure 7, we observe that for vector mesons, $D^*$, at temperature $T = 0$ and baryon density $\rho_0 = \rho_0 (4\rho_0)$, the values of mass shift due to condensates $\langle \bar{q} q \rangle$, $\langle q^4 D_0 q^4 \rangle$, $\langle q_g \sigma G q \rangle$, $\langle q_g D_0 q \rangle$, $\langle \alpha G G / \pi \rangle$, and $\langle \alpha G G / \pi \rangle$ are observed to be $-82.15 (-151.15)$, $1.689 (6.749)$, $18.95 (57.16)$, $-5.49 (-6.16)$, and $3.54 (14.93)$ MeV, respectively. For temperature $T = 100$ MeV the above values of mass shift at baryon density $\rho_0 = \rho_0 (4\rho_0)$ changes to $-73.01 (-142.35)$, $1.689 (6.75)$, $17.85 (56.34)$, $-4.53 (-5.43)$, and $3.43 (14.64)$ MeV, respectively. For the axial-vector mesons, $B_1$, the values of mass shift due to individual condensates $\langle \bar{q} q \rangle$, $\langle q^4 D_0 q^4 \rangle$, $\langle q_g \sigma G q \rangle$, $\langle q_g D_0 q \rangle$, and $\langle \alpha G G / \pi \rangle$ are observed to be $79.31 (152)$, $0.1516 (0.607)$, $-8.77 (-22.63)$, $8.27 (21.76)$, and $2.34 (10.08)$ MeV, respectively. For temperature $T = 100$ MeV the above values of mass shift at baryon density $\rho_0 = \rho_0 (4\rho_0)$ changes to $70 (145)$, $0.151 (0.607)$, $-7.99 (-22.04)$, $7.58 (21.26)$, and $2.23 (9.81)$, respectively. In Figures 9, 10, and 11 we have shown the contributions of individual condensates to the mass shift of scalar mesons, $B_0$, vector mesons, $B^*$, and axial-vector mesons, $B_1$, respectively. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. In Tables 1, 2, and 3 we have tabulated the values of mass shift for scalar mesons, $B_0$, vector mesons, $B^*$ and axial-vector mesons, $B_1$, respectively. The values of mass shift have been given at baryon densities $\rho_0$ and $4\rho_0$ and temperatures $T = 0$ and 100 MeV.

In [23], the properties of scalar mesons $D_0$ and $B_0$ had been studied using QCD sum rules and the observed values of mass shift at nuclear saturation density were 69 and 217 MeV, respectively. The properties of $D_1$ mesons had also been studied in [25] using coupled channel approach. An extra widening from the already large width of the resonance in free space was observed for the $D_1$ meson. The properties of vector mesons ($D^*$ and $B^*$) and axial-vector ($D_1$ and $B_1$) had been studied using QCD sum rules in [24]. For vector mesons $D^*$ and $B^*$ the values of mass shift were $-71$ and $-380$ MeV, respectively, whereas for axial-vector mesons $D_1$ and $B_1$ these values change to 72 and 264 MeV, respectively. The observed positive values of mass shift for $D_0$ and $D_1$ mesons in our present investigation and also in earlier investigations indicate that these mesons feel repulsive interactions in the nuclear medium and their in-medium mass increases as function of baryonic density. This means the chances of decay of higher charmonium states to these heavy charmed mesons pairs are suppressed and hence these mesons may not cause a decrease in the production of $J/\psi$ mesons in heavy-ion collisions.

### Table 1: The following table gives the values of mass shift of scalar mesons $B_0$ (in MeV units) due to individual condensates at temperatures $T = 0$ and 100 MeV. For each value of temperature the values are tabulated for baryon densities $\rho_0 = \rho_0$ and $4\rho_0$.

| $T$ (MeV) | $\rho_0 = \rho_0$ | $\rho_0 = 4\rho_0$ |
|----------|-----------------|-----------------|
| $T = 0$  | 224             | 196             | 399             |
| $T = 100$| 223             | 195             | 402             |

| $T$ (MeV) | $\rho_0 = \rho_0$ | $\rho_0 = 4\rho_0$ |
|----------|-----------------|-----------------|
| $T = 0$  | -367            | -318.08         | -665.88         |
| $T = 100$| -373.11         | 1.477           | 5.9057          |
| $T = 0$  | 12.43           | 11.67           | 36.18           |
| $T = 100$| -3.98           | -6.42           | -3.40           |
| $T = 0$  | 1.64            | 7.56            | 1.51            |
| $T = 100$| -366.34         | -307.26         | -317.28         |

### Table 3: (Color online) The following table gives the values of mass shift of axial vector mesons $B_1$ (in MeV units) due to individual condensates at temperatures $T = 0$ and 100 MeV. For each value of temperature the values are tabulated for baryon densities $\rho_0 = \rho_0$ and $4\rho_0$.

| $T$ (MeV) | $\rho_0 = \rho_0$ | $\rho_0 = 4\rho_0$ |
|----------|-----------------|-----------------|
| $T = 0$  | 266.51          | 232.75          | 467.29          |
| $T = 100$| 270.11          | 236.35          | 478.86          |
| $T = 0$  | 0.88            | 0.88            | 3.53            |
| $T = 100$| -6.34           | -5.83           | -16.40          |
| $T = 0$  | 4.79            | 4.39            | 12.23           |
| $T = 100$| 1.41            | 6.39            | 6.12            |
| $T = 0$  | 266.75          | 493.11          | 232.99          | 468.20          |
Figure 9: (Color online) Figure shows the contribution of individual condensates to the mass shift of scalar vector mesons $B_0$ as a function of the squared Borel mass parameter, $M^2$. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. 
Figure 10: Figure shows the contribution of individual condensates to the mass shift of vector mesons $B^*$ as a function of the squared Borel mass parameter, $M^2$. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. 
Figure 11: (Color online) Figure shows the contribution of individual condensates to the mass shift of axial-vector mesons $B_1$ as a function of the squared Borel mass parameter, $M^2$. The subplots (a), (c), and (e) show the results at temperature $T = 0$, whereas the subplots (b), (d), and (f) are plotted for temperature $T = 100$ MeV. We have shown the results at baryon densities $\rho_0$, $2\rho_0$, and $4\rho_0$. 
collisions. However, for the vector mesons $D^*$ we observe the negative values of mass shift and hence we conclude that they feel attractive interactions in the nuclear medium. The decrease in the mass of vector mesons $D^*$ may cause the decay of higher charmonium states to $D^0\overline{D}^*$ pairs and hence it may be a cause of $J/\psi$ suppression in heavy-ion collision experiments.

5. Summary

In the present paper we investigated the mass modifications of scalar mesons ($D_0, B_0$), vector mesons ($D^*, B^*$) and axial-vector mesons ($D_1, B_1$) at finite density and temperature of the nuclear medium. We used QCD sum rules along with chiral SU(3) model to investigate the properties of above mentioned mesons. Using chiral SU(3) model we found the values of quark and gluon condensates at finite density and temperature of the medium which were further used within QCD sum rules to find the in-medium masses of scalar, vector, and axial-vector mesons. We observed a positive value of mass shift for the scalar mesons ($D_0, B_0$) and axial-vector mesons ($D_1, B_1$); that is, their in-medium mass was found to be more than the vacuum mass. For a constant value of temperature the values of mass shift for these scalar and axial-vector mesons are found to increase as we increase the density of the nuclear medium. On the other hand for a constant value of density the temperature of the medium causes a decrease in the values of mass shift of scalar and axial-vector mesons. The vector mesons $D^*$ and $B^*$ are found to have negative values of mass shift. It means their in-medium masses are small as compared to vacuum masses. For a constant value of density, as a function of temperature, the magnitude of mass shift of $D^*$ and $B^*$ mesons decreases. However, for a constant value of temperature, as a function of density the magnitude of mass shift of vector mesons $D^*$ and $B^*$ are found to increase. We have also investigated the effects of individual terms on the mass shift of mesons. It was found that the scalar quark condensates, $\bar{q}q$, have maximum contribution to the in-medium modification of scalar, vector, and axial-vector mesons. The effects of density and temperature of the medium on the scattering lengths of scalar, vector, and axial-vector mesons were also investigated. We observed the positive value of scattering lengths for scalar and axial-vector mesons, whereas for the vector mesons a negative value of scattering lengths were observed. Also it was found that the magnitude of the scattering lengths of scalar, vector, and axial-vector mesons decreases as we move from low to high value of density or temperature of the nuclear medium. The present investigation of medium modification of scalar, vector, and axial-vector mesons will be helpful for understanding their production rate and also the phenomenon of $J/\psi$ suppression in the compressed baryonic matter experiment at FAIR, GSI.

Conflict of Interests

The author declares that there is no conflict of interests regarding the publication of this paper.

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