Photon-Hadron Jet Correlations in $p+p$ and $Au+Au$ Collisions at $\sqrt{s_{NN}}=200$ GeV

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We report the observation at the Relativistic Heavy Ion Collider (RHIC) of suppression of back-to-back correlations in the direct photon+jet channel in Au+Au relative to p+p collisions. Two-particle correlations of direct photon triggers with associated hadrons are obtained by statistical subtraction in perturbative quantum chromodynamics (pQCD), making such correlations a powerful probe of the in-medium parton energy loss. The away-side nuclear suppression factor, $I_{AA}$, in central Au+Au collisions, is $0.32 \pm 0.12^{\text{syst}} \pm 0.09^{\text{stat}}$ for hadrons of 3 $< p_T < 5$ in coincidence with photons of 5 $< p_T < 15$ GeV/c. The suppression is comparable to that observed for high-$p_T$ single hadrons and dihadrons. The direct photon associated yields in $p+p$ collisions scale approximately with the momentum balance, $z_T \equiv p_T^{\gamma} / p_T^{h}$, as expected for a measure of the away-side parton fragmentation function. We compare to Au+Au collisions for which the momentum balance dependence of the nuclear modification should be sensitive to the path-length dependence of parton energy loss.

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I. INTRODUCTION

Experimental results from RHIC have established the formation of hot and dense matter of a fundamentally new nature in relativistic heavy-ion collisions at $\sqrt{s_{NN}} = 200$ GeV [1]. Energy loss in this dense nuclear matter by color-charged, hard (E $\gtrsim$ 2 GeV) partons, and the jets into which they fragment, is generally accepted to be the mechanism responsible for the suppression of the high-$p_T$ hadron yields observed in central A+A collisions [2,3]. In the large multiplicity environment of heavy-ion collisions, two-particle correlations are often used to study jet modification and to infer properties of the medium. For example, high-$p_T$ azimuthal dihadron correlations demonstrate that the degree of jet away-side suppression depends on the $p_T$ of the “trigger” and “associated” hadrons. At moderate $p_T$ ($\gtrsim$ 3 GeV/c), the jet properties measured through two-particle correlations demonstrate novel features such as shape modifications which are thought to be a manifestation of the response of medium to the energy deposited by the attenuated parton [4].

Di-hadron measurements of dijet pairs provide an ambiguous measurement of the energy loss of the away-side parton. The trigger hadron is a product of parton fragmentation and therefore it is not possible to determine, event-by-event, whether the near-side parton has itself lost energy. Given the steeply falling jet spectrum, the sample of hard scatterings is biased towards configurations in which the parton loses little energy. In particular, it is believed that hadron measurements are subject to a “surface bias” in which the hard scatterings sampled are likely to occur at the periphery of the overlap zone [3,6]. The away-side parton then is more likely to traverse a maximal path-length through the medium. For a sufficiently opaque medium, the attenuation of the parton may be nearly total, in which case the sensitivity to the average path-length is reduced [7]. Back-to-back, high-$p_T$ hadron pairs may originate preferentially from configurations in which the outgoing parton trajectories are tangential to the surface of the overlap zone [8]. On the other hand, dihadron pairs may also originate from vertices deep in the collision zone if a parton has a finite probability to “punch-through” or pass through the medium without interaction [9]. Calculations of the relative importance of these two mechanisms depend both on the model of parton energy loss employed and the density

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profile of the medium \[6,10,11\].

Direct photon-jet pairs offer two major advantages in studying energy loss as compared to dijets because of the nature of the photon. First, in contrast to partons, photons do not carry color charge and hence do not interact strongly when traversing the medium \[12\]. The distribution of hard scattering vertices sampled by direct photon-triggered correlations is thus unbiased by the trigger condition. Suppression of the opposite jet is averaged over all path-lengths given by the distribution of hard scattering vertices. Second, at the Born level, direct photon-triggered correlations is thus unbiased by the trigger condition. Suppression of the opposite jet is averaged over all path-lengths given by the distribution of hard scattering vertices. In this way, the average path-length of the away-side parton may then be varied in a well controlled manner by selecting events of various momentum differences between the \(\gamma\)-h pair.

For this reason, the \(\gamma\)-jet channel has long been considered the “golden channel” for studying parton energy loss \[13,14\]. Neglecting the above mentioned complications, specifically effects like transverse momentum broadening (the \(k_T\) effect) and parton-to-photon fragmentation, back-to-back \(\gamma\)-h correlations in elementary collisions directly measure the fragmentation function of the recoil jet since \(z = p^h/p^{\mathrm{jet}} \approx p^h/p^\gamma\). In the standard picture of energy loss, partons are likely to lose some fraction of their energy in the medium, but are likely to fragment outside the medium. Hence, the parton energy loss can be considered an effective modification to the fragmentation function. Such a picture may be tested using \(\gamma\)-h correlations in nuclear collisions. Complementary baseline measurements in \(p+p\) collisions are used to test the theoretical description of correlations in vacuum and to constrain possible contributions from higher order processes. Comprehensive reviews of direct photon phenomenology and data from elementary collisions may be found in \[15,16,17\].

II. DETECTOR DESCRIPTION AND PARTICLE IDENTIFICATION

The data were taken with the PHENIX detector \[18\] using approximately 950 million \(Au+Au\) minimum bias events from the 2004 data set and 471 million photon-triggered events from the 2005 and 2006 \(p+p\) data sets corresponding to integrated luminosities of 3 \(\text{PB}(2005)\) and 10.7 \(\text{PB}(2006)\). The Beam-Beam Counters (BBC) \[19\], which are used to trigger the minimum bias data, select 92% of the total inelastic cross section. In \(Au+Au\) the BBC and Zero-Degree Calorimeters (ZDC) were used for offline minimum bias event selection and centrality determination. In \(p+p\) collisions a high energy photon trigger, defined by coincidence between the BBC and a high energy Electromagnetic Calorimeter (EMCal) tower hit, was utilized. This EMCal based trigger \[20\] had an efficiency of \(> 90\%\) for events with photons and \(\pi^0\) with energies in the range used in the analysis and within the detector’s geometric acceptance.

The PHENIX central arms, each covering \(\pm 0.35\) units of pseudorapidity around midrapidity and \(90^\circ\) in azimuth, contain charged-particle tracking chambers and electromagnetic calorimeters \[21\]. The EMCal \[22\] consists of two types of detectors, six sectors of lead-scintillator (PbSc) sampling calorimeters and two of lead-glass (PbGl) Čerenkov calorimeters measuring EM energy with intrinsic resolution \(\sigma_E/E = 8.1%/\sqrt{E} \oplus 2.1\%\) and \(5.9%/\sqrt{E} \oplus 0.8\%\) respectively. The fine segmentation of the EMCal \((\Delta p \times \Delta \phi) \sim 0.01 \times 0.01\) for PbSc and \(\sim 0.008 \times 0.008\) for PbGl allows for the reconstruction of \(\pi^0\) and \(\eta\) mesons in the \(2\gamma\) decay channel out to \(p_T\) of 20 GeV/c. The details of direct photon, \(\pi^0\) and \(\eta\) meson detection and reconstruction within PHENIX have been described previously \[12,22,24\]. Photon candidates with very high purity \((> 98\%\) for energies > 5 GeV) are selected from EMCal clusters with the use of cluster shower shape and charged particle veto cuts. Two-photon \(\pi^0\) and \(\eta\) candidates are selected from photon pairs with pair invariant mass in the appropriate \(\pi^0\) or \(\eta\) mass range. Combinatorial \(2\gamma\) background is reduced with cuts on energy asymmetry \(\alpha_{12} = |E_1 - E_2|/(E_1 + E_2)\), described in detail below. Some fraction of \(\pi^0\) with \(p_T\) starting at \(\sim 13\) GeV/c (in the PbSc detector) will appear as a single merged cluster, but with anomalous shower shape, and thus are removed from the analysis. The \(\pi^0\) and \(\eta\) mesons in the \(p_T\) range from about 4 to 17 GeV/c and photons between 5 and 15 GeV/c are used in this analysis. For \(\gamma\) \(p_T\) between 13–15 GeV/c there is a \(< 2\%\) contribution of merged \(\pi^0\) cluster contamination, however this together with all sources of non-photon contamination are found to have a negligible impact on the two-particle correlation analysis of this report. Direct photons and their two-particle correlations are obtained by statistical subtraction of the estimated meson (mainly \(\pi^0\)) decay photon contribution from the inclusive photon and \(\gamma\)-h samples.

Charged hadrons are detected with the PHENIX tracking system \[27\] which employs a drift chamber in each arm spanning a radial distance of 2.0–2.4 m from the beam axis with a set of pixel pad chambers (PC1) directly behind them. The momentum resolution was determined to be \(\delta p/p = 0.7% \oplus 1.0\%p\) where \(p\) is measured in GeV/c. Secondary tracks from decays and conversions are suppressed by matching tracks to hits in a second pad chamber (PC3) at distances of \(\sim 5.0\) m. Track projections to the EMCal plane are used to veto photon candidates resulting from charged hadrons that shower in the EMCal.
III. METHOD

A. Two-Particle Correlations

Two-particle correlations are constructed by measuring the yield of particle pairs as a function of the measured azimuthal angle between photon or parent meson triggers and charged hadron partners. The correlation function, \( C(\Delta\phi) \equiv N^\text{pair}(\Delta\phi)/N^\text{mixed}(\Delta\phi) \), corrects for the limited acceptance of \( \gamma-h \) or meson-hadron pairs by dividing the distribution in real events \( N^\text{pair} \) by the mixed event distribution \( N^\text{mixed} \). The correlation function is decomposed utilizing a two-source model of pair yields coming from two-particle jet correlations superimposed on a combinatorial background yield from an underlying event. The underlying event in Au+Au is known to have an azimuthal asymmetry of harmonic shape quantified in the elliptic flow parameter \( v_2 \) [26, 27]. This flow represents a harmonic modulation of the \( \Delta\phi \) distribution of this underlying event, such that the flow-subtracted jet correlation signal is encoded in the jet pair ratio function, \( JPR(\Delta\phi) \equiv C(\Delta\phi) - \xi(1 + 2\langle v_2^2 \rangle \cos 2\Delta\phi) \), using the notation of [1], where \( \langle v_2 \rangle \) is the average single-particle \( v_2 \).

Two methods of determining the background level \( \xi \), known as Zero-Yield at Minimum (ZYAM) and Absolute Normalization (ABS) respectively were applied to the Au+Au data. Both methods are described in detail in previous PHENIX publications [4], see also [26, 29] (ABS) and [30] (ZYAM). ZYAM assigns the level of zero jet yield and hence \( \xi \) to the minimum point of the correlation function \( C(\Delta\phi) \). The ABS method uses the mean multiplicity of trigger-associated pairs in mixed events and a correction for finite centrality resolution to determine \( \xi \). Where ZYAM statistical precision is reasonable, the direct \( \gamma-h \) extraction of the two methods agree to within much better than the total uncertainties, typically within \( \lesssim 20\% \). The ABS method is chosen for the Au+Au results presented, as this method resulted in a more precise extraction of direct photon-jet pair yields at high trigger \( p_T \) where lack of statistics near \( \Delta\phi = \pi/2 \) severely impairs the ZYAM determination. In the comparatively low multiplicity \( p+p \) collisions, the underlying event originates from different physical mechanisms than in Au+Au and is known not to be well described by event-mixing. Instead the correlation functions are normalized by fitting to a double Gaussian + constant function, corresponding to the ZYAM method [4].

The results presented here are corrected for the associated charged hadron efficiency \( \epsilon_h \) such that the quoted yields correspond to a detector with full azimuthal acceptance and \( |\eta| < 0.35 \) coverage. No correction is applied for the \( \Delta\eta \) acceptance of pairs. Final results are presented in terms of the yield \( Y \) of jet pairs per trigger, \( Y \equiv A JPR(\Delta\phi)/N_{\text{trigger}} \) with the constant \( A = \int N^\text{pair}(\Delta\phi)/(2\pi\epsilon_h) \).

The magnitudes of elliptic flow were determined by measuring the distributions of inclusive photons, neutral pions, and charged hadrons as a function of the angle relative to the reaction plane, which was determined with the BBC’s as described in [31]. The \( v_2 \) values measured for this analysis are consistent with previous PHENIX analyses [21, 27, 32]. At high-\( p_T \) (\( \geq 6 \text{ GeV}/c \)) the measured \( \pi^0 \) \( v_2 \) values used in the determination of the decay photon \( v_2 \) are fit to a constant function in order to reduce the effects of large statistical fluctuations. The \( p_T \) independence of \( v_2 \) of \( \pi^0 \) is motivated by recent preliminary data [33] and also by the observed \( p_T \) independence of the \( R_{AA} \), since parton energy loss is expected to be the dominant mechanism for \( v_2 \) generation at high-\( p_T \) [7]. It is also consistent with the findings of [32] which is direct measurement of \( \pi^0 \) \( v_2 \) for the same dataset and is being published concurrently with this measurement. Since, as discussed in that publication, the high-\( p_T \) functional behavior for this dataset cannot be well-constrained, the level of uncertainty we assign to the constant fit assumption increases with \( p_T \).

Table I lists the \( v_2 \) values for the inclusive and \( \pi^0 \) decay photons for all \( p_T \) ranges used, either the measurements, or for the highest \( p_T \) decay \( v_2 \) values from the constant fit value. For the fit values the fit errors are listed as statistical error, despite the inherent systematic correlation of the fit value across the \( p_T \) bins. The decay photon \( v_2 \) is derived from the measured \( \pi^0 \) \( v_2 \) by the same \( p_T^{\pi^0} \rightarrow p_T^{\text{decay}\gamma} \) mapping procedure applied to the yields, described below. It is assumed that the \( v_2 \) for other mesons which contribute decay photons (e.g. \( \eta \)) are the same as that of the \( \pi^0 \) at high-\( p_T \). This assumption is well motivated for the \( p_T \) range considered (\( \gtrsim 4.5 \text{ GeV}/c \)) under the expectation that the source of the high \( p_T \) azimuthal asymmetry \( v_2 \) is jet quenching-induced suppression, already measured to be the same for a variety of mesons (e.g. \( \eta \) itself [24]) and by data measurements for other high \( p_T \) \( v_2 \) which confirm the expectation [34] for other hadrons.

| Centrality | \( p_T^{\pi^0} \) | \( v_2 \) Stat. | Sys. | \( v_2 \) Stat. | Sys. |
|------------|------------------|------------------|------------------|------------------|------------------|
| 5-7        | 0.053 ± 0.009    | 0.011            | 0.084 ± 0.009    | 0.004            |
| 7-9        | 0.047 ± 0.022    | 0.015            | 0.069 ± 0.018    | 0.003            |
| 9-12       | 0.024 ± 0.017    | 0.0063           | 0.069 ± 0.020    | 0.003            |
| 12-15      | 0.004 ± 0.096    | 0.094            | 0.069 ± 0.023    | 0.003            |
| 5-7        | 0.096 ± 0.010    | 0.005            | 0.155 ± 0.011    | 0.036            |
| 7-9        | 0.079 ± 0.027    | 0.011            | 0.105 ± 0.019    | 0.025            |
| 9-12       | 0.025 ± 0.050    | 0.049            | 0.105 ± 0.020    | 0.025            |
| 12-15      | 0.287 ± 0.128    | 0.104            | 0.105 ± 0.023    | 0.024            |
| 5-7        | 0.143 ± 0.023    | 0.035            | 0.136 ± 0.022    | 0.010            |
| 7-9        | 0.146 ± 0.064    | 0.026            | 0.126 ± 0.039    | 0.008            |
| 40-60%     | 9-12             | 0.162 ± 0.126    | 0.252            | 0.126 ± 0.042    | 0.008            |
| 12-15      | -0.603 ± 0.308   | -0.191           | 0.126 ± 0.046    | 0.008            |
TABLE II: Extracted \( R_\gamma \) values used as input to direct \( \gamma-h \) per-trigger yield subtraction (Equation 2). These values are interpolated from previous PHENIX measurements as described in the text.

| Centrality | \( p_T^\gamma \) (GeV/c) | \( R_\gamma \) | Stat. | Sys. |
|------------|-------------------------|----------------|-------|------|
| 5-7        | 1.77                    | ±0.09          | ±0.06 |
| 7-9        | 2.45                    | ±0.09          | ±0.18 |
| 0-20%      | 2.99                    | ±0.11          | ±0.41 |
| 12-15      | 3.66                    | ±0.24          | ±0.68 |
| 20-40%     | 1.46                    | ±0.10          | ±0.04 |
| 7-9        | 1.85                    | ±0.10          | ±0.12 |
| 9-12       | 2.30                    | ±0.12          | ±0.28 |
| 12-15      | 2.35                    | ±0.20          | ±0.44 |
| 40-60%     | 1.30                    | ±0.09          | ±0.05 |
| 7-9        | 1.52                    | ±0.07          | ±0.13 |
| 9-12       | 1.85                    | ±0.10          | ±0.30 |
| 12-15      | 1.94                    | ±0.24          | ±0.36 |
| \( p+p \)  | 5-7                     | 1.18           | ±0.01 | ±0.06 |
| 7-9        | 1.33                    | ±0.01          | ±0.05 |
| 9-12       | 1.53                    | ±0.03          | ±0.05 |
| 12-15      | 1.79                    | ±0.09          | ±0.07 |

The per-trigger yield of inclusive \( \gamma-h \) pairs \( Y_{\text{inclusive}} \) is simply the weighted average of the contributions from decay and direct photon triggers,

\[
Y_{\text{inclusive}} = \frac{N^\gamma_{\text{direct}} Y_{\text{direct}} + N^\gamma_{\text{decay}} Y_{\text{decay}}}{N_{\text{inclusive}}}, \tag{1}
\]

Having already determined \( R_\gamma \), \( Y_{\text{direct}} \) may then be obtained by simple manipulation of the above terms resulting in statistical subtraction involving only per-trigger yields as follows. The decay photon per-trigger yield is subtracted from that of inclusive photons according to:

\[
Y_{\text{direct}} = \frac{R_\gamma Y_{\text{inclusive}} - Y_{\text{decay}}}{R_\gamma - 1} \tag{2}
\]

The direct \( \gamma \) or direct \( \gamma-h \) pair yields do not, by definition, exclude photons from jet fragmentation or medium induced photon production.

C. Extraction of Decay Photon Correlations

The decay photon associated yields are estimated from the measured \( \pi^0-h \) and \( \eta-h \) correlations through a calculation which determines the decay correlations statistically from a Monte Carlo (MC) based, pair-by-pair weighting procedure. In this procedure the decay \( \gamma-h \) pair yield \( N^\gamma_{\text{decay}}(p_T^\gamma) \) is constructed by a weighted integral over all \( \pi^0-h \) and \( \eta-h \) pairs. In what follows, we will first describe the procedure schematically, describing the ingredients and how they are obtained. We then give a more exact description and associated formula representing exactly how the weighting was performed in the measurement. Schematically the procedure may be expressed as a convolution of several factors according to the following relation, wherein for simplicity we only consider photons from \( \pi^0 \) decay, although the procedure is also applied to \( \eta \) decay photons.

\[
N^\gamma_{\text{decay}}(p_T^\gamma) = \frac{\epsilon_\pi(p_T^\pi) \otimes \mathcal{P}(p_T^\pi)}{\epsilon_\pi(p_T^\pi)} \otimes N^{\pi^0-h}(p_T^\pi) \tag{3}
\]

where \( \epsilon_\pi \) and \( \gamma \) are the \( \pi^0 \) and single decay photon efficiencies, respectively, and \( \mathcal{P} \) is the decay probability density, each of which is addressed in turn below.

First, since the starting point is the uncorrected raw meson-h pair yield \( N^{\pi^0-h} \), a correction for the parent meson reconstruction efficiency, \( \epsilon_\pi(p_T^\pi) \), is applied to the raw \( \pi^0 \)’s as a function of \( p_T \) in order to account for the \( \pi^0 \) daughters in the inclusive sample whose sisters lie outside the PHENIX acceptance or are otherwise undetected. Both efficiencies \( \epsilon_\pi \) and \( \epsilon_\gamma \) in Equation 3 are also evaluated as a function of the position in the calorimeter along the beam direction, however this dependence mostly cancels in the ratio \( \epsilon_\gamma/\epsilon_\pi \) and therefore is suppressed for clarity. \( \epsilon_\pi(p_T^\pi) \) is determined by dividing the raw number of \( \pi^0 \)’s \( N^{\pi^0-h}(p_T^\pi) \) obtained in the same data sample by the PHENIX published \( \pi^0 \) invariant yields [2, 24, 38] assuming no pseudorapidity dependence over the narrow PHENIX acceptance. The top panel in Fig. 1 illustrates, for the example of central Au+Au events, the \( \pi^0 \) efficiency correction factor \( 1/\epsilon_\pi(p_T^\pi) \). The correction rises at small \( p_T \) due to a \( p_T \)-dependent pair energy asymmetry cut designed to reduce combinatorial 2-\( \gamma \) pairs reconstructed as real \( \pi^0 \)’s. This cut, along with
Analytic smearing of the resolution. Occupancy effects give rise to an additional contribution according to the known EMCal energy and position Gaussian smearing functions to simulate detector resolution. A simple fast MC generator implements the PHENIX acceptance and uses the single decay photon efficiency, $\epsilon_\gamma(p_T^\gamma)$, which differs from the true decay photon distribution. The single decay photon efficiency depends on both the parent and daughter $p_T$ and is evaluated in a GEANT simulation. In principle the convolution of both $P(p_T^\gamma, p_T^\pi)$ and $\epsilon_\gamma(p_T^\gamma, p_T^\pi)$, $W(p_T^\gamma, p_T^\pi)$, could be extracted as one function from the GEANT simulation, but obtaining large enough MC statistics necessary to properly parameterize the above mentioned EMCal $z$ position dependence of the $\epsilon_\gamma$ corrections is only feasible with the fast MC. Thus only the efficiency loss by cluster merging for photons $\epsilon_{\gamma,merge}$ is taken from the GEANT. The bottom panel of Fig. 1 shows $\epsilon_{\gamma,merge}(p_T^\gamma)$ evaluated from the GEANT simulation.

Since we wish to construct per-trigger yields, the same procedure described in Equation (3) can be applied to find the estimated single decay photon trigger yield from the measured single $\pi^0$s, i.e. replacing $N_{\text{decay}}$ with $N_{\text{decay}}^{\pi^0}$ and $N_{\pi^-h}$ with $N_{\pi^-}$. The exact application of schematic Equation (3) then takes the form of a sum over all $p_T^0, h$ pairs and single $\pi^0$'s found in the data. Each $\pi^0$ or $\pi^-h$ pair is given a weight which depends on $p_T^0$. Operationally we now split this weight into two parts: $\epsilon_\gamma(p_T^\gamma)$ discussed above and a factor $W_{ab}(p_T^\gamma)$. The factor $W_{ab}$ is simply the end result of the fast MC-GEANT combined calculation, the convolution of $P$ and $\epsilon_\gamma$, including $\epsilon_{\gamma,merge}$, averaged over a chosen decay photon bin of the range $a < p_T < b$. Thus in terms of the product $W(p_T^\gamma, p_T^\pi)$ then $W_{ab}(p_T^\gamma)$ is given by

$$W_{ab}(p_T^\gamma) = \int_a^b dp_T^\pi W(p_T^\gamma, p_T^\pi)$$

Terms $W_{ab}(p_T^\gamma)$ are defined for the four photon $p_T$ bins used in the analysis, $[a, b] = [5, 7], [7, 9], [9, 12]$ and $[12, 15]$ GeV/c. An example of $W_{ab}(p_T^\gamma)$ for the 5-7 GeV/c bin is shown in Fig. 1. Procedurally, we construct $W_{ab}$ as product of the fast MC curve shown in the middle panel and the linear fit discussed above to the bottom panel, $\epsilon_{\gamma,merge}(p_T^\gamma)$. Although a decay of $p_T^\gamma < a$, the lower limit of the decay $p_T$ bin, is kinematically disallowed, $W_{ab}$ is non-zero below this boundary when resolution effects are considered. For $p_T^\gamma > b$, $W_{ab}$ decreases as $\sim 1/p_T^\gamma$, slowly enough that $\pi^0$'s at values of $p_T$ beyond the statistical reach of the data set contribute to the relevant decay photon $p_T$ selections at a non-negligible rate. The $\pi^0$ sample is truncated at $p_T = 17$ GeV/c and extrapolated using power-law fits to the single and conditional $\pi^0$ spectra to estimate a correction. In the latter case, each associated hadron $p_T$ range is fit independently. The truncation avoids the high-$p_T$ region where cluster merging the effects of any remaining background, is described below. At large $p_T$, $1/\epsilon_\gamma(p_T^\gamma)$ rises again due to losses from cluster merging.

Second, the effect of decay kinematics is evaluated by determining the probability density, $P(p_T^\gamma, p_T^\pi)$, for the decay of a $p_T$-independent distribution of $\pi^0$’s. $P(p_T^\gamma, p_T^\pi)$ represents the relative probability of a $\pi^0$ of $p_T = p_T^\gamma$, to decay into a photon of $p_T^\pi$. For a perfect detector, this function is calculable analytically. A simple fast MC generator implements the PHENIX acceptance and uses Gaussian smearing functions to simulate detector resolution according to the known EMCal energy and position resolution. Occupancy effects give rise to an additional smearing of the $\pi^0$ and $\eta$ invariant masses. This effect is included in the MC by tuning the resolution parameters to match the $\pi^0$ peak widths observed in data. False reconstruction of $\pi^0$’s and $\eta$’s from combinatorial matches are either subtracted or assigned to the systematic uncertainties as discussed below.

Finally, we wish to estimate the decay photon contribution to the measured raw inclusive photon sample which differs from the true decay photon distribution by the single decay photon efficiency, $\epsilon_\gamma(p_T^\gamma)$. At intermediate $p_T$, $\epsilon_\gamma(p_T^\gamma)$ depends only on the photon momentum and is included already implicitly by the fast MC simulation described above to produce $P(p_T^\gamma, p_T^\pi)$. Thus, it is useful to think of them as a single factor $W(p_T^\gamma, p_T^\pi) \equiv P(p_T^\gamma, p_T^\pi) \epsilon_\gamma(p_T^\gamma, p_T^\pi)$. At high-$p_T$, on the other hand, an efficiency loss is incurred by photons from $\pi^0$’s whose showers merge into a single cluster in the calorimeter and are rejected by the shower-shape cut. As a consequence, the fraction of photons that are direct is artificially enhanced in the sample of reconstructed photon clusters. The single decay photon efficiency depends on both the parent and daughter $p_T$ and is evaluated in a GEANT simulation. In principle the convolution of both $P(p_T^\gamma, p_T^\pi)$ and $\epsilon_\gamma(p_T^\gamma, p_T^\pi)$, $W(p_T^\gamma, p_T^\pi)$, could be extracted as one function from the GEANT simulation, but obtaining large enough MC statistics necessary to properly parameterize the above mentioned EMCal $z$ position dependence of the $\epsilon_\gamma$ corrections is only feasible with the fast MC. Thus only the efficiency loss by cluster merging for photons $\epsilon_{\gamma,merge}$ is taken from the GEANT. The bottom panel of Fig. 1 shows $\epsilon_{\gamma,merge}(p_T^\gamma)$ evaluated from the GEANT simulation.

FIG. 1: The weight factors used to obtain decay correlations from parent meson correlations. Top: $\pi^0$ reconstruction efficiency correction, $1/\epsilon_\gamma$. Middle: Decay probability function, $W_{ab}$, for 5–7 GeV/c decay photons from $\pi^0$ derived analytically (black line), using the detector acceptance and resolution smearing (red line) and including the single decay photon efficiency, $\epsilon_\gamma$, from a GEANT simulation (blue points). Bottom: $\epsilon_{\gamma,merge}$ obtained by taking ratio of the blue points to red curve in the previous panel.
effects are dominant and the $1/\epsilon^\pi$ correction factor becomes large. Although the truncation corrections for the number of decay photons and decay $\gamma$-$h$ pairs are non-negligible, they mostly cancel in the per-trigger yield and are therefore typically < 1%, reaching a maximum value of 7% for only the $12 < p_T^\gamma < 15 \otimes 3 < p_T^h < 5$ GeV/c bin.

With the weight functions $W_{ab}$ the entire set of $\pi^0$-hadron pairs and single $\pi^0$ candidates (within a given range of $\Delta\phi$, $\phi_1 < \Delta\phi < \phi_2$, defining each $\Delta\phi$ bin) are then summed over, once for each decay photon $p_T^\gamma$ bin, and the per-trigger yield is constructed for each of these decay $p_T$ bins as

$$Y_{\text{decay}}(\phi_1 < \Delta\phi < \phi_2) = \frac{\sum_{i=1-N}^{N} W_{ab}(p_T^{\gamma i})/\epsilon_\pi(p_T^{\gamma i})}{\sum_{i=1-N}^{N} W_{ab}(p_T^{h i})/\epsilon_h(p_T^{h i})}$$ (5)

In this form it is clear that the normalization of the functions $\epsilon_\pi(p_T^{\gamma})$ and $W_{ab}(p_T^{\gamma})$ cancel out completely in the per-trigger yield, and therefore only their shapes versus $p_T^{\gamma}$ are important. Hence in Fig. 1 the curves are shown with arbitrary units. Also, as Equation 5 implies, the angular deviation between the direction of a decay photon and its parent meson is ignored. The $\Delta\phi$ opening angle of a decay photon and hadron pair is taken to be the same as the $\Delta\phi_{\pi-h}$ of the parent $\pi^0$-$h$ pairs. This approximation is tested in the fast MC and found to be extremely accurate since the distribution of angular deviation between a leading decay photon in a 2$\gamma$ decay and the parent mesons at these $\pi^0$ momenta have an RMS around 0 of $\approx 0.01$ radians, and the smallest $\Delta\phi$ bins considered in the analysis are typically $\sim 0.1$ radians or larger.

D. $\pi^0$ and $\eta$ Reconstruction

In $p+p$ collisions $Y_{\text{decay}}$ is estimated using both reconstructed $\pi^0$ and $\eta$ mesons in invariant mass windows of 120–160 and 530–580 MeV/c$^2$, respectively. The total decay per-trigger yield is calculated from

$$Y_{\text{decay}} = (1 - \delta_{h/\pi^0}) Y_{\text{decay}}^\pi + \delta_{h/\pi^0} Y_{\text{decay}}^\eta$$ (6)

where $\delta_{h/\pi^0}$ is the ratio of the total number of decay photons to the total number of decay photons from $\pi^0$. Based on the measurements of $\eta$ $[24]$ and $\omega$ $[33]$, which together with the $\pi^0$ account for > 99% of decay photons, the value of $\delta_{h/\pi^0}$ is determined to be 1.24±0.05 in the high-$p_T$ region covered by this analysis, independent of collision system and centrality. Note that the per-trigger yields for $\omega$ and other heavier meson triggers ($\omega, \eta, \phi, ...$) are not measured and are taken to be equivalent to $Y_{\text{decay}}^{\pi^0}$ in Equation 6. This assumption was studied in PYTHIA and found to influence $Y_{\text{decay}}$ at the level of < 2%. In Au+Au collisions correlations using $\eta$ triggers are not directly measured, but rather estimated from the $p+p$ measurement as discussed below.

Figure 2 shows the various components of the decay photon measurement in $p+p$. In $p+p$ collisions the rate of combinatorial background photon pairs is reduced by only considering photons of $p_T > 1$ GeV/c resulting in background levels of < 10% for which no correction was applied. The effect of such remaining pairs on $Y_{\text{decay}}^\pi$ was evaluated to be negligible (< 2%) compared to the size of other uncertainties on the final result using a detailed full PYTHIA test of the method which included $\pi^0$ reconstruction with combinatorial photon pairs. On the other hand, $\eta$ reconstruction has a much smaller signal-to-background of 1.4–1.6, depending on the $p_T$ selection, even in the low multiplicity $p+p$ environment. In this case, the per-trigger yield of the combinatorial photon pairs is estimated from photon pairs with invariant mass in “sideband” ranges of 400–460 and 640–700 MeV/c$^2$, beyond 3$\sigma$ of the $\eta$ peak. The sideband contribution $Y_{\text{decay}}^{\text{sideband}}$ is then subtracted using the signal-to-background ratio $f_{\text{bkg}}$ evaluated from gaussian + polynomial background fits to the invariant mass distributions according to $Y_{\text{signal}}^{\text{decay}} = Y_{\text{signal}}^{\text{raw}}/(1/f_{\text{bkg}} + 1) - Y_{\text{sideband}}^{\text{decay}}/f_{\text{bkg}}$. The yield $Y_{\text{sideband}}^{\text{decay}}$ is generated from the full meson to decay photon weighting function procedure (Equation 5). The subtraction procedure was also tested in PYTHIA and the extracted and input per-trigger yields were found to agree to within 10%.

In Au+Au collisions the combinatorial rate for $\pi^0$ reconstruction is substantially larger. Correspondingly, a $p_T$ dependent cut on the pair energy asymmetry $a_{12} = |E_1 - E_2|/(E_1 + E_2)$ $[23]$, visible in Fig. 1 with the small-

![FIG. 2: (color online) Examples of parent and daughter per-trigger yields for the $\pi^0$ and $\eta$ in $p+p$ collisions for $p_T$ selection $5 < p_T^\gamma < 7$ and $2 < p_T^h < 3$ GeV/c. These correlation measurements are used to determine the total decay photon per-trigger yield as described in the text.](image-url)
est allowed asymmetry at the lowest $\pi^0$ $p_T$ values, is used to reduce this background. With such cuts the signal-to-background in central events varies from 5:1 at its lowest, increasing to about 15:1 for the highest $p_T$ selection. The effect of the combinatorial background is studied through examination of a similar sideband subtraction analysis as in the $p+p$ $\eta-h$ correlation extraction described, this time for $\pi^0-h$, using invariant mass ranges just outside the $\pi^0$ peak region. However no clear trend beyond non-negligible statistical limitations is observed, so no correction for the background is applied. Instead the maximum size of the effect (typically $\sim 7\%$) is included as source of systematic uncertainty on the decay yields and propagated to the final direct photon per-trigger yields.

In central Au+Au collisions the $\eta$ meson cannot be reconstructed with sufficient purity to measure its correlations. Instead, a scaling argument is employed. Motivated by the similar high-$p_T$ suppression pattern shown by $\eta$ and $\pi^0$ in Au+Au and corresponding near equality of the $p+p$ and Au+Au $\eta/\pi^0$ ratios, the ratio $Y_{\gamma}(\eta)/Y_{\gamma}(\pi^0)$ is measured in $p+p$ and applied as a correction to the Au+Au $Y_{\gamma}(\pi^0)$. This is justified by the assumption that the jet fragmentation is primarily occurring outside the medium. We do not attribute any additional uncertainty to this scaling beyond the 10% sideband systematic and statistical uncertainties of the $\eta$ measurement in $p+p$ yields. However, to give an idea of the possible impact of this assumption, the total systematic error on $Y_{\gamma}(\eta)$ from all other sources would correspond to a variation of the Au+Au $Y_{\gamma}(\pi^0)$ by $\sim 50\%$. Given the similarity of the high-$p_T$ suppression demonstrated by all light quark bound states measured thus far, this would correspond to a rather large change.

IV. SYSTEMATIC UNCERTAINTIES

There are four main classes of systematic uncertainty in the Au+Au data: elliptic flow, normalization of the underlying event (ABS), $R_\gamma$, and the decay per-trigger yield estimate, the latter two of which are present in the $p+p$ data as well. Table III lists the fractional contribution of each of these sources to the total systematic uncertainty on the direct photon per-trigger yields in the 20% most central Au+Au and $p+p$ data. In the central Au+Au data the uncertainty at low $p_T^h$ is dominated by the $v_2$ and correlation function normalization (ABS method) estimation due to large multiplicity of hadrons. At higher $p_T^h$, but low trigger $p_T$, $p_T^h$, the decay error dominates due to the two-photon combinatorial background for $\pi^0$ reconstruction. Finally, at large $p_T^h$ the backgrounds responsible for both of these sources of uncertainty decrease and the uncertainty on $R_\gamma$, which is relatively constant, dominates. In $p+p$ collisions the decay photon background forms a much larger fraction of the total photon sample. In this case, the decay uncertainty arises from the MC decay photon mapping procedure, the $\eta$ sideband subtraction and the $\eta/\pi^0$ ratio in approximately equal parts. The yields associated with daughter photons are larger than for the meson parents because of feed-down from larger values of parent $p_T$, and hence, jet $p_T$.

The correction for single hadron efficiency $c_h(p_T^h)$ varies as a function of collision system and centrality. These corrections are obtained by finding the ratio of raw yields of hadrons obtained without the trigger constraints to the previous PHENIX published measurements of the corresponding charged hadron spectra [40, 41]. As in previous PHENIX two-particle correlation measurements [4, 30], this procedure has inherent uncertainties assigned as a $p_T$-independent 10% uncertainty, on each system and/or centrality.

V. RESULTS

A. Direct $\gamma-h$ Per-Trigger Yields

Figure 3 shows examples of direct photon per-trigger yields in $p+p$ and central Au+Au collisions. Also shown are the per-trigger yields for inclusive and decay photon triggers which are the ingredients in the statistical subtraction method as expressed in Equation 2. A clear away-side correlation is observed ($\Delta\phi \simeq \pi$) for direct photons triggers in $p+p$. In Au+Au collisions the away-side correlation is suppressed for both decay and direct photon triggers. The near-side direct photon associated yields are small relative to that of decay photons, an expected signature of prompt photon production [16].

The away-side yields, integrated over $|\Delta\phi - \pi| < \pi/5$ radians, are shown in Fig. 3 and Table IV for $p+p$ and Au+Au collisions. This range roughly corresponds to the “head region” as defined in 4 and is chosen primarily to minimize the influence of medium response which is thought to dominate the “shoulder” region further offset from $\Delta\phi = \pi$. Additionally, the acceptance and the
Inclusive $\gamma$-h Decay

FIG. 4: (color online) Direct $\gamma$-h per-trigger yields for the range $|\Delta\phi - \pi| < \pi/5$ radians vs. associated hadron $p_T$. Four different direct $\gamma$ $p_T$ ranges (indicated on the figure) are shown in the most central 20% of Au+Au events and $p+p$ events. The upper limits are for 90% confidence levels. A $p_T$-independent uncertainty of 10% due to the charged hadron efficiency corrections is not shown.

signal itself are largest in this range so statistical precision is maximized. It should be noted that the width of the jet correlation is larger than this interval. We do not make a correction for this effect, since we are primarily concerned with the comparison of the yields from $p+p$ and Au+Au collisions. It should be noted, however, that in addition to parton energy loss, any broadening of azimuthal correlations, whether by hot or cold nuclear matter effects, will contribute to a suppression in the yield in the head region. Due to statistical and systematic fluctuations, the subtraction of the decay-photon hadron pairs from the inclusive $\gamma$-h sample can result in a negative yield. In this case 90% confidence-level upper limits are given. In the case that a positive yield is obtained, but the uncertainty is consistent with 0, the lower bound of the error bar is also replaced with an arrow. As noted in the figure caption, a 10% $p_T$-independent uncertainty due to the charged hadron efficiency corrections is not shown.

**TABLE IV: Direct $\gamma$-h per-trigger yields in 20% most central Au+Au and in $p+p$ collisions. An additional $p_T$-independent uncertainty of 10% due to the charged hadron efficiency corrections is not shown.**

| $p_T^\gamma$ (GeV) | $p_T^h$ (GeV) | Yield (Stat) | Sys | Total |
|-------------------|---------------|--------------|-----|-------|
| Au+Au, Centrality 0–20% |
| 1-2 | 0.23 | 6.26e-02 | 4.62e-02 | 6.60e-02 |
| 5-7 | 0.41 | 2.68e-02 | 5.68e-03 | 1.41e-02 |
| 3-5 | 0.62 | 4.82e-03 | 2.13e-03 | 2.90e-03 |
| 7-9 | 0.17 | 3.71e-02 | 5.59e-02 | 1.02e-01 |
| 9-12 | 0.3 | 3.45e-02 | 2.39e-03 | 2.53e-02 |

**B. Suppression Factor $I_{AA}$**

Departure from the vacuum QCD processes is quantified by $I_{AA}$, the ratio of Au+Au to $p+p$ per-trigger yields:

$$I_{AA}(p_T^\gamma, p_T^h) = \frac{Y^{Au+Au}(p_T^\gamma, p_T^h)}{Y^{p+p}(p_T^\gamma, p_T^h)}.$$  \hspace{1cm} (7)
resulted in a negative yield value (the 90% confidence level) for associated hadron bins for the most central 0–20% of collisions. The data points for which the subtraction is not shown. Figure 5 shows the ratio of suppression plays an important role as would be expected from a sample dominated by surface emission. 

two measurements should be subject to different geometrical effects. Disentangling such effects through precise comparisons of dihadron and $\gamma$-h suppression should be pursued with future measurements with improved statistics.

C. Towards the Fragmentation Function

Using the distribution of charged hadrons opposite direct $\gamma$ triggers, parton energy loss may be studied directly as a departure from the (vacuum) fragmentation function. In distinction to $\pi^0$-h correlations, where the away-side distribution is only sensitive to the integral of the fragmentation function (the average multiplicity of the away-side jet), the away-side distribution for direct $\gamma$-h correlations provides a measurement of the full fragmentation function. In distinction to $pT$-h correlations, the away-side distribution is consistent with a scenario in which the geometric of suppression plays an important role as would be expected from a sample dominated by surface emission.

Figure 5 also compares $I_{AA}$ from a measurement of high-$p_T$ dihadron ($h^+ - h^\pm$) correlations to the $\gamma$-jet result for similar $p_{T,h}$ selections. The two results are remarkably similar in the most central bin. This may indicate that surface emission is dominant for both samples in this $z$ region. However it should be noted that the total uncertainties on either measurement are still quite large on a relative scale. As explained in the introduction, the

FIG. 5: (color online) Ratio $I_{AA}$ of the Au+Au to $p+p$ yields shown in Fig. 4. An additional $p_T$-independent uncertainty of 14% due to the charged hadron efficiency corrections is not shown.

FIG. 6: (color online) $I_{AA}(p_T^\gamma)$ integrated over the range $5 < p_T^\gamma < 15$ GeV/c for associated hadrons of $3 < p_T^h < 5$ GeV/c vs. centrality compared to single $\pi^0$-high-$p_T$: $R_{AA}$ (integrated over $p_T > 5$ GeV/c). An additional $p_T$-independent uncertainty of 14% due to the charged hadron efficiency corrections is not shown.
away parton do not exactly balance. The transverse momentum imbalance was discovered at the CERN-ISL using $x_E$ distributions [43] and originally attributed to an “intrinsic” transverse momentum $k_T$ of each of the initial colliding partons [44], but now understood to be due to “resummation” of soft-gluon effects [45, 46].

The validity of the approximation $x_E \approx z_\gamma$ can be tested by observing identical $x_E$ distributions for different values of trigger $p_T^\gamma$ ($x_E$ scaling), in which case one would accept the $x_E$ distribution in $\gamma$-$h$ correlations as the quark fragmentation function from the reaction $q+g \rightarrow q+\gamma$ without need of correction. We approximate $x_E$ by $z_\gamma$, the ratio of the mean associated $p_T^h$ to mean trigger $p_T^\gamma$ for each $p_T^\gamma$ bin. The $(p_T^h)$ for the four trigger bins are: 5.66, 7.75, 10.07, 13.07 GeV/c, close to the values obtained from a fit to the direct-$\gamma$ invariant cross section of the form $p_T^{-b.5}$.

Figure 7 shows the $z_T$ distributions for p+p and Au+Au collisions. The p+p data (Fig. 7a) exhibit reasonable $z_T$ scaling so that the measured distribution should represent the away-side jet fragmentation function. A fit of this data to a simple exponential ($N e^{-b z_T}$) gives an acceptable $\chi^2/\text{dof} = 12.8/10$ with a value $b = 6.9 \pm 0.8$ which is consistent with the quark fragmentation function, parameterized [42] as a simple exponential with $b = 8.2$ for $0.2 < z < 1.0$, and inconsistent with the gluon fragmentation function value of $b = 11.4$. It should, however, be recalled that the data do not cover the full extent of the away peak, only $|\Delta \phi - \pi| < \pi/5$ radians, and that possible variations of the widths of the peaks in both the p+p data and the Au+Au data with $p_T^h$ and $p_T^\gamma$ have not been taken into account in the present analysis. Additionally a more detailed analysis, differential in trigger $p_T$, is necessary to study trigger $p_T$ dependent effects which can influence the fragmentation function fit values [42].

In central Au+Au collisions, the fragmentation function may be modified by the medium\(^2\), so that $z_T$ scaling should not hold except in two special cases: i) pure surface emission or punch-through where the away-side jets are not modified—the $z_T$ distribution will be suppressed, but will have the same shape as in p+p collisions; ii) constant fractional energy loss of the away jet—the $z_T$ scaling will be preserved in Au+Au collisions but with a steeper slope than in p+p collisions. The Au+Au data (Fig. 7b) are consistent with $z_T$ scaling with the same shape as the $p+p$ data, a value of $b = 5.6 \pm 2.2$ and excellent $\chi^2/\text{dof} = 10.1/10$ for the simple exponential fit. The point at lowest $z_T = 0.11$ for Au+Au is 1.6 standard deviations above the fit, suggesting that improved statistics will permit the observation of any non-surface emission.

D. Model Comparison

Several authors have reported predictions for $\gamma$-jet in heavy ion collisions [48, 49, 50, 51]. As a demonstration of the how such calculations can be compared to the data, the $I_{AA}$ values as a function of $z_T$ are compared to energy loss predictions [49] in Fig. 8. The calculation uses effective fragmentation functions to parameterize the energy

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1 The reader is advised to carefully distinguish this variable $z_T = \langle p_T^h \rangle / \langle p_T^\gamma \rangle$ from our previous notation used in [42] of $z$ = $p_T^h / \bar{p}$, which is the fraction of jet momentum $\bar{p}$ contained in the trigger particle.

2 See Equation 1 in [47]
loss in terms of a parameter $\varepsilon_0$ which is expected to be proportional to the initial gluon density $\rho_0$. The model calculates the energy-loss of the leading parton, and neglects the contribution the gluon radiation and medium response which may dominate at low values of $z$. The data is well reproduced by the model over the range of values of $\varepsilon_0$ provided, 1.48–1.88 GeV/fm. This corresponds roughly to the range of $\varepsilon_0$ allowed by comparison to the PHENIX $\pi^0 R_{AA}$ data of $1.5^{+0.2}_{-0.5}$. It should be noted that the calculation rejects fragmentation photons with an isolation cut. Such a procedure has not yet been demonstrated in central Au+Au data, although doing so would help to eliminate beyond-leading-order effects.

VI. CONCLUSIONS

We have presented the first direct $\gamma$-h measurements in Au+Au and $p+p$ collisions at RHIC. A significant suppression of $I_{AA} = 0.32 \pm 0.12^{\text{stat}} \pm 0.09^{\text{syst}}$ for the away-side charged hadron yield in the range $3 < p_T < 5$ GeV/c is observed for direct photon triggers in Au+Au as compared to $p+p$. Furthermore, the level of suppression is found to be consistent with the single particle suppression rate and the importance of energy-loss geometry, notably the expectation of surface emission in the kinematic range sampled. A possible indication that energy-loss geometry may also be important in dijet suppression is that $\gamma$-h suppression $I_{AA}$ is also observed to be quite similar to that of dihadron suppression in central events; however, the current precision of the data does not exclude substantial differences. In the $p+p$ data $z_T$ scaling is observed, suggesting that the measured $z_T$ distribution (Fig. 6) is a statistically acceptable representation of the fragmentation function of the quark jet recoiling away from the direct photon. Improvement of the statistical and systematic precision of the measurements should allow further tests of vacuum fragmentation expectations in $p+p$ collisions and insights into details of the medium modification of jet fragmentation in Au+Au.

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