Symmetries of Lorentzian Three-Manifolds with Recurrent Curvature

Giovanni CALVARUSO † and Amirhesam ZAEIM ‡

† Dipartimento di Matematica e Fisica “E. De Giorgi”, Università del Salento, Prov. Lecce-Arnesano, 73100 Lecce, Italy
E-mail: giovanni.calvaruso@unisalento.it
‡ Department of Mathematics, Payame Noor University, P.O. Box 19395-3697, Tehran, Iran
E-mail: zaeim@pnu.ac.ir

Received February 12, 2016, in final form June 17, 2016; Published online June 26, 2016
http://dx.doi.org/10.3842/SIGMA.2016.063

Abstract. Locally homogeneous Lorentzian three-manifolds with recurrent curvature are special examples of Walker manifolds, that is, they admit a parallel null vector field. We obtain a full classification of the symmetries of these spaces, with particular regard to symmetries related to their curvature: Ricci and matter collineations, curvature and Weyl collineations. Several results are given for the broader class of three-dimensional Walker manifolds.

Key words: Walker manifolds; Killing vector fields; affine vector fields; Ricci collineations; curvature and Weyl collineations; matter collineations

2010 Mathematics Subject Classification: 53C50; 53B30

1 Introduction

A Walker manifold is a pseudo-Riemannian manifold \((M, g)\) admitting a degenerate parallel distribution. Such a phenomenon is peculiar to the case of indefinite metrics. As such, it is responsible for many special geometric properties of pseudo-Riemannian manifolds which do not have any Riemannian counterpart, and has been investigated by several authors under different points of view. The monograph [5] is a well-written recent survey on Walker manifolds and the various related research areas.

Lorentzian three-manifolds admitting a parallel degenerate line field have been studied in [12]. These Lorentzian metrics are described in terms of a suitable system of local coordinates \((t, x, y)\) and form a large class, depending on an arbitrary function \(f(t, x, y)\). The case of strictly Walker manifolds, where the parallel degenerate line field is spanned by a parallel null vector field, is characterized by condition \(f = f(x, y)\). The results of [12] have been recently used in [14] to obtain a complete classification of the models of locally homogeneous Lorentzian three-manifolds with recurrent curvature.

The aim of this paper is to investigate symmetries of these Lorentzian spaces. If \((M, g)\) denotes a Lorentzian manifold and \(T\) a tensor on \((M, g)\), codifying some either mathematical or physical quantity, a symmetry of \(T\) is a one-parameter group of diffeomorphisms of \((M, g)\), leaving \(T\) invariant. As such, it corresponds to a vector field \(X\) satisfying \(\mathcal{L}_X T = 0\), where \(\mathcal{L}\) denotes the Lie derivative. Isometries are a well known example of symmetries, for which \(T = g\) is the metric tensor. The corresponding vector field \(X\) is then a Killing vector field. Homotheties and conformal motions on \((M, g)\) are again examples of symmetries. In recent years, symmetries related to the curvature of the manifold have been investigated. Among them: curvature collineations (where \(T = R\) is the curvature tensor), Weyl collineations (\(T = W\) being the Weyl conformal curvature tensor) and Ricci collineations, for which \(T = \varrho\) is the Ricci tensor. We
may refer to the monograph [16] for further information and references on symmetries. Ricci and curvature collineations have been investigated in several classes of Lorentzian manifolds (see, for example, [1, 6, 7, 8, 9, 13, 15, 17, 18, 19, 20, 21, 22] and references therein). Because of their physical relevance, in most cases curvature symmetries have been studied for some space-times. Moreover, the three-dimensional case has also been considered as an interesting source of examples and different behaviours (see, for example, [8]).

A \textit{matter collineation} of a Lorentzian manifold \((M,g)\) is a vector field \(X\), corresponding to a symmetry of the energy-momentum tensor \(T = \varrho - \frac{1}{2} \tau g\), where \(\tau\) denotes the scalar curvature. Matter collineations are more relevant from a physical point of view [10, 11], while Ricci collineations have a more clear geometrical significance, since \(\varrho\) is naturally deduced from the connection of the metric [19]. These physical and geometrical meanings do coincide in a special case, namely, for metrics with vanishing scalar curvature. And this is exactly the case for any strictly Walker three-manifold [12].

We shall obtain complete classifications of curvature and Ricci (≡ matter) collineations of homogeneous Lorentzian three-manifolds with recurrent curvature. In Section 2 we shall give some basic information about Walker three-manifolds and curvature symmetries. In Section 3 we then investigate symmetries of an arbitrary strictly Walker three-manifold. Since the function \(f = f(x,y)\) determining the metric tensor here is arbitrary, one cannot expect to obtain these symmetries explicitly in the general case. However, we describe the sets of partial differential equations describing the different symmetries and use them to give some explicit examples of proper symmetries. Then, in Section 4 we shall completely classify the symmetries of homogeneous Lorentzian three-manifolds with recurrent curvature. All calculations have also been checked using \textit{Maple16}. 

2 Preliminaries

2.1 Three-dimensional Walker metrics

We shall essentially follow the notations used in [12]. A three-dimensional Lorentzian manifold \(M\) admitting a parallel degenerate line field has local coordinates \((t,x,y)\), such that with respect to the local frame field \(\{\partial_t, \partial_x, \partial_y\}\) the Lorentzian metric is given by

\[
g_f = \begin{pmatrix}
0 & 0 & 1 \\
0 & \varepsilon & 0 \\
1 & 0 & f(t,x,y)
\end{pmatrix},
\]

for some function \(f(t,x,y)\). In the above expression, \(\varepsilon = \pm 1\). However, it is easily seen that by reversing the metric and changing the sign of the coordinate \(x\), without loss of generality one can reduce to the case \(\varepsilon = 1\) (as it was done, for example, in [14]).

The parallel degenerate line field is spanned by \(\partial_t\), and the existence of a parallel null vector \(U = \partial_t\) (\textit{strictly Walker metric}) is characterized by the independence of the function \(f\) of the variable \(t\) [23]. Therefore, with respect to local coordinates \((t,x,y)\), the general form of a strictly Walker metric is given by

\[
g_f = \begin{pmatrix}
0 & 0 & 1 \\
0 & 1 & 0 \\
1 & 0 & f(x,y)
\end{pmatrix},
\]  

(2.1)

for an arbitrary smooth function \(f\). With respect to the coordinate basis \(\{\partial_t, \partial_x, \partial_y\}\), the Levi-Civita connection \(\nabla\) and curvature \(R\) of the metric \(g_f\) described by (2.1) are completely determined by the following possibly non-vanishing components (see also [12]):

\[
\nabla_{\partial_x} \partial_y = \frac{1}{2} f_x \partial_t, \quad \nabla_{\partial_y} \partial_y = \frac{1}{2} f_y \partial_t - \frac{1}{2} f_x \partial_x
\]  

(2.2)
and
\[ R(\partial_x, \partial_y)\partial_x = \frac{1}{2} f_{xx} \partial_t, \quad R(\partial_x, \partial_y)\partial_y = -\frac{1}{2} f_{xx} \partial_x, \tag{2.3} \]
where \( R(X, Y) = [\nabla_X, \nabla_Y] - \nabla_{[X,Y]} \). From (2.2) and (2.3), a straightforward calculation yields that the covariant derivative of the curvature tensor is completely determined by the possibly non-vanishing components
\[
(\nabla_{\partial_x} R)(\partial_x, \partial_y)\partial_x = \frac{1}{2} f_{xx} \partial_t, \quad (\nabla_{\partial_x} R)(\partial_x, \partial_y)\partial_y = -\frac{1}{2} f_{xx} \partial_x, \\
(\nabla_{\partial_y} R)(\partial_x, \partial_y)\partial_x = \frac{1}{2} f_{xy} \partial_t, \quad (\nabla_{\partial_y} R)(\partial_x, \partial_y)\partial_y = -\frac{1}{2} f_{xy} \partial_x. \tag{2.4} \]

Either by (2.4) or by direct calculation differentiating the Ricci identity, it is easily seen that three-dimensional (strictly) Walker metrics have \textit{recurrent curvature}, that is, in a neighborhood of any point of non-vanishing curvature, one has \( \nabla R = \omega \otimes R \), for a suitable one-form \( \omega \). Since we are interested in the study of the nonflat examples with recurrent curvature, throughout the paper we shall assume that \( f_{xx} \neq 0 \) at any point.

In local coordinates \((t, x, y)\), the Ricci tensor \( \rho \) of any metric (2.1) is given by
\[
\rho = \begin{pmatrix}
0 & 0 & 0 \\
0 & 0 & 0 \\
0 & -\frac{1}{2} f_{xx}
\end{pmatrix}. \tag{2.5}
\]

A pseudo-Riemannian manifold \((M, g)\) is said to be \textit{locally homogeneous} if for any pair of points \( p, q \in M \) there exist a neighbourhood \( U \) of \( p \), a neighbourhood \( V \) of \( q \) and an isometry \( \phi: U \to V \). Hence, locally homogeneous manifolds “look the same” around each point. For any given class of pseudo-Riemannian manifolds, it is a natural problem to determine its locally homogeneous examples. Locally homogeneous examples among three-dimensional Walker metrics have been investigated in [14] (see also [3]). Rewriting the classification obtained in [14] in terms of coordinates \((t, x, y)\) used in (2.1), we have the following.

\textbf{Theorem 2.1 ([14])}. Locally homogeneous Lorentzian three-manifolds of recurrent curvature naturally divide into three classes. They correspond to one of the following types of (strictly) Walker metrics, as described in (2.1):

\begin{enumerate}
  \item [I)] \( N_b \), defined by taking \( f(x, y) = -2b^{-2} e^{bx} \), for some real constant \( b \neq 0 \);
  \item [II)] \( P_c \), defined by taking \( f(x, y) = -x^2 \alpha(y) \), where \( \alpha > 0 \) satisfies \( \alpha_y' = c \alpha^{3/2} \) for some real constant \( c \);
  \item [III)] \( CW_\varepsilon \), defined by taking \( f(x, y) = -\varepsilon x^2 \), where \( \varepsilon = \pm 1 \).
\end{enumerate}

\section{2.2 Curvature and Ricci Collineations}

Let \((M, g)\) denote a pseudo-Riemannian manifold (in particular, a Lorentzian one). A vector field \( X \) on \( M \) preserving its metric tensor \( g \), the corresponding Levi-Civita connection \( \nabla \), its curvature tensor \( R \) or its Ricci tensor \( \rho \), is respectively known as a \textit{Killing vector field}, an \textit{affine vector field}, a \textit{curvature collineation} or a \textit{Ricci collineation}.

It is obvious that if \( X \) preserves \( g \) (respectively, \( \nabla, R \)), then it also preserves \( \nabla \) (respectively, \( R, \rho \)), but the converse does not hold in general. Homothetic vector fields (i.e., vector fields \( X \) satisfying \( \mathcal{L}_X g = \lambda g \) for some real constant \( \lambda \)) are again necessarily curvature collineations (in particular, Ricci collineations). For this reason, we are specifically interested in the existence of \textit{proper} Ricci and curvature collineations, namely, the ones which are not homothetic (and hence, not Killing). Thus, we also need to specify which are the Killing, affine and homothetic vector
fields, which is an interesting problem on its own, due to the natural geometric meaning of such symmetries.

Conditions defining Ricci and curvature collineations are formally similar to the ones defining Killing or affine vector fields. However, they may show some deeply different behaviours. In fact (see, for example, [16, 19]):

(a) Killing and affine vector fields are smooth (provided they are at least $C^1$). However, for any positive integer $k$, there exist Lorentzian metrics admitting Ricci (and curvature) collineations, which are $C^k$ but not $C^{k+1}$.

(b) Unlike Killing and affine vector fields, Ricci (and curvature) collineations form a vector space which may be infinite-dimensional and (because of the above point (a)) is not necessarily a Lie algebra. In fact, if $X, Y$ are Ricci (curvature) collineations, then $[X, Y]$ might not be differentiable.

(c) While Killing and affine vector fields agreeing in the neighbourhood of a point must coincide everywhere, two Ricci (respectively, curvature) collineations that agree on a non-empty subset of $M$ may not agree on $M$, since they are not uniquely determined by the value of $X$ and its covariant derivatives of any order at a point.

Observe that the above item (b), as concerns the possibility of the vector space of Ricci collineations to be infinite-dimensional, refers to cases where the Ricci tensor $\rho$ is necessarily degenerate (as it is always the case, for example, for three-dimensional strictly Walker metrics). On the other hand, if $\rho$ (respectively, $T = \rho - \frac{1}{2} \tau g$) is nondegenerate, then Ricci (respectively, matter) collineations form a finite-dimensional Lie algebra of smooth vectors. In fact, in such a case, they are exactly the Killing vector fields of the nondegenerate metric tensor $\rho$.

3 Symmetries of Walker three-manifolds

Observe that any three-dimensional strictly Walker metric is already equipped in a natural way with the strongest possible symmetry. In fact, it admits a parallel vector field $U = \partial_t$. We shall now investigate the different kinds of symmetries of these metrics, starting with Killing, homothetic and affine vector fields. The descriptions we obtain are given in the following.

**Theorem 3.1.** Let $(M, g_f)$ be a three-dimensional strictly Walker manifold, where $g_f$ is described in the local coordinates $(t, x, y)$ by equation (2.1). A smooth vector field $X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y$ is

i) a Killing vector field if and only if

$$X_1 = -c_1 t - x f_1'(y) + f_2(y), \quad X_2 = f_1(y), \quad X_3 = c_1 y + c_2,$$

where $f_1, f_2$ are smooth functions on $M$, satisfying

$$2c_1 f - 2f_1''(y)x + 2f_2'(y) + f_1(y)\partial_x f + (c_1 y + c_2)\partial_y f = 0.$$ (3.2)

ii) a homothetic, non-Killing vector field if and only if

$$X_1 = \eta t - c_1 t - x f_1'(y) + f_2(y), \quad X_2 = \frac{\eta}{2} x + f_1(y), \quad X_3 = c_1 y + c_2,$$

where $\eta \neq 0$ is a real constant and

$$(2c_1 - \eta) f - 2f_1''(y)x + 2f_2'(y) + \left(\frac{\eta}{2} x + f_1(y)\right)\partial_x f + (c_1 y + c_2)\partial_y f = 0.$$ (3.4)
iii) an affine Killing vector field if and only if

\[ X_1 = c_3 t - x f'_1(y) + f_2(y), \quad X_2 = \frac{c_1 + c_3}{2} x + f_1(y), \quad X_3 = c_1 y + c_2, \]  

(3.5)

where

\[
(c_1 - c_3) f - 2 f''_1(y) x + 2 f'_2(y) + \left( \frac{c_1 + c_3}{2} x + f_1(y) \right) \partial_x f
\]

\[
+ \left( c_1 y + c_2 \right) \partial_y f + c_4 = 0.
\]

(3.6)

**Proof.** We start from an arbitrary smooth vector field \( X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y \) on the three-dimensional strict Walker manifold \((M, g_f)\), where \( g_f \) is described by equation (2.1), and calculate \( \mathcal{L}_X g_f \). Then, \( X \) satisfies \( \mathcal{L}_X g_f = \eta g_f \) for some real constant \( \eta \) if and only if the following system of partial differential equations is satisfied:

\[
\partial_t X_3 = 0, \quad \partial_x X_2 = \frac{\eta}{2}, \quad \partial_x X_3 + \partial_t X_2 = 0, \quad \partial_y X_3 + \partial_t X_3 f + \partial_t X_1 = \eta,
\]

\[
f \partial_x X_3 + \partial_y X_1 + \partial_y X_2 = 0, \quad 2 \partial_y X_3 f + 2 \partial_y X_1 + X_2 \partial_x f + X_3 \partial_y f = \eta f.
\]

(3.7)

We then proceed to integrate (3.7). From the first three equations in (3.7) we get \( X_2 = \frac{\eta}{2} x - a_1(y) t + f_1(y) \) and \( X_3 = a_1(y) x + b_1(y) \). Then, the fourth equation in (3.7) yields \( X_1 = \eta t - a'_1(y) t x - b'_1(y) t + f_4(x, y) \). Substituting this into the fifth equation, we get

\[
2a'_1(y) t = a_1(y) f + \partial_x f_4(x, y) + f'_2(y),
\]

which must hold for all values of \( t \), implying that \( a_1(y) = c_1 \) is a constant. Now, the last equation in (3.7) gives

\[
(c_1 \partial_x f + 2 b'_1(y)) t = (2 b'_1(y) - \eta f) f + 2 \partial_y f_4(x, y) + \left( \frac{\eta}{2} x + f_2(y) \right) \partial_x f
\]

\[
+ \left( c_1 x + f_3(y) \right) \partial_y f,
\]

which immediately yields that \( c_1 \partial_x f + 2 b'_1(y) = 0 \) and so, \( c_1 \partial_x^2 f = 0 \). Since we assumed \( \partial_x^2 f \neq 0 \), we then have \( c_1 = 0 \) and integrating \( b''_1(y) = 0 \) we get \( b_1(y) = c_2 y + c_3 \). On the other hand, from the fifth equation in (3.7) we now have \( f_4(x, y) = -f'_2(y) x + f_5(y) \) and the last equation gives

\[
(2 c_2 - \eta) f - 2 f'_2(y) x + 2 f'_3(y) + \left( \frac{\eta}{2} x + f_2(y) \right) \partial_x f + (c_2 y + c_3) \partial_y f = 0.
\]

This proves the statement i) in the case \( \eta = 0 \) and the statement ii) if we assume \( \eta \neq 0 \).

With regard to affine vector fields, expressing condition \( \mathcal{L}_X \nabla = 0 \) in the coordinate basis \( \{ \partial_t, \partial_x, \partial_y \} \), we get the following system of partial differential equations:

\[
\partial_{tt} X_1 = \partial_{tt}^2 X_2 = \partial_{tx}^2 X_2 = \partial_{tx} X_2, \quad \partial_{xx}^2 X_3 = \partial_{xx} X_3 = \partial_{tx}^2 X_3 = \partial_{ty}^2 X_3 = 0,
\]

\[
\partial_{tx}^2 X_1 + \partial_x X_3 \partial_x f = 0,
\]

\[
2 \partial_{tx}^2 X_1 + \partial_x X_3 \partial_x f = 0,
\]

\[
2 \partial_{ty} X_2 - \partial_x X_3 \partial_x f = 0,
\]

\[
2 \partial_{xy}^2 X_3 - \partial_x X_3 \partial_x f = 0,
\]

\[
2 \partial_{xy}^2 X_1 + \partial_x X_3 \partial_y f + \partial_t X_2 \partial_x f = 0,
\]

\[
2 \partial_{xy}^2 X_2 - \partial_x X_3 \partial_x f - \partial_t X_2 \partial_x f = 0,
\]

\[
2 \partial_{yy}^2 X_3 + \partial_x X_3 \partial_x f - \partial_t X_3 \partial_y f = 0,
\]
2\partial_y X_3 \partial_x f - 2\partial_y^2 X_2 - 2\partial_y X_2 \partial_y f + \partial_y X_3 \partial_x f + X_2 \partial_y^2 f + X_3 \partial_y^2 f = 0,
\partial_y X_3 \partial_y f + 2\partial_y^2 X_1 + \partial_y X_3 \partial_y f + \partial_y X_2 \partial_y f - \partial_y X_1 \partial_x f + X_2 \partial_y^2 f + X_3 \partial_y^2 f = 0,
2\partial_y X_3 \partial_y f + 2\partial_y^2 X_1 + \partial_y X_1 \partial_x f + 2\partial_y X_2 \partial_x f - \partial_y X_1 \partial_y f + X_2 \partial_y^2 f + X_3 \partial_y^2 f = 0. \quad (3.8)

As for the above system (3.7), we then proceed to integrate (3.8). From the first equation we get $X_3 = c_1 t + a_1(y) x + f_2(y)$ and then the fifth equation yields $2a_1'(y) = c_1 \partial_x f$, so that $c_1 \partial_x^2 f = 0$ and so, $c_1 = 0$. Then, $a_1(y) = c_2$ is a constant.

Integrating the third and fourth equations (taking into account the first one) we get $X_1 = f_3(y)t + f_4(x, y)$, $X_2 = c_3 t + f_5(y)x + f_6(y)$. The sixth equation then gives $f_3'(y) + c_3 \partial_x f = 0$, which, by the same argument above, yields $c_3 = 0$ and $f_3(y) = c_4$.

Similarly, the eighth equation becomes $2f_7''(y) + c_2 \partial_x f = 0$, which gives $c_2 = 0$ and $f_2(y) = c_5 y + c_6$. The second equation now reads $\partial_x^2 f_4(x, y) = 0$, and the seventh leads to $f_5'(y) = 0$.

So, $f_4(x, y) = f_7(y)x + f_8(y)$ and $f_5(y) = c_7$.

By the ninth and tenth equations we then have $f_9''(y) + f_7'(y) + \frac{1}{2}(2c_7 - c_4 - c_5) \partial_x f = 0$, so that $c_7 = \frac{c_4 + c_5}{2}$ and $f_7(y) = -f_9'(y) + c_8$. Integrating the tenth equation with respect to the variable $x$, we get

$$
(c_5 - c_4) f - 2f_9''(y)x + \left(\frac{c_4 + c_5}{2} x + f_6(y)\right) \partial_x f + (c_5 y + c_6) \partial_y f + f_9(y) = 0.
$$

We differentiate the above equation with respect to $y$ and subtract the eleventh equation, obtaining $f_9'(y) - 2f_9''(y) = c_8 \partial_x f$, which immediately leads to $c_8 = 0$ and $f_9(y) = 2f_8'(y) + c_9$. The statement follows after we suitably rename the remaining constants and functions.

**Remark 3.2.** Since $f = f(x, y)$ is an arbitrary smooth function of two variables, we cannot integrate equations (3.2), (3.4) and (3.6) of Theorem 3.1 in full generality. However, it is well known that the Lie algebras of Killing, homothetic and affine vector fields are finite-dimensional.

Therefore, on the one hand, for any prescribed function $f(x, y)$ these equations force the sets of Killing, homothetic and affine vector fields of $(M, g_f)$ to depend on a finite number of real parameters. On the other hand, they allow us to determine special functions $f$, for which we can find some explicit examples of homothetic non-Killing and affine non-homothetic vector fields.

At the end of this section we shall illustrate these results calculating the symmetries of an arbitrary locally conformally flat strictly Walker metric. In the next section we shall consider the functions $f$ determining the locally homogeneous examples of three-dimensional strictly Walker manifolds. Further explicit examples can be determined by direct calculation.

**Remark 3.3** (homothetic fixed points). The existence on a Lorentzian manifold $(M, g)$ of homothetic fixed points, that is, of a non-trivial homothetic vector field $X$ which vanishes at a point $m \in M$, has some important consequences on the structure of the manifold itself. Different conclusions can be deduced depending on whether $m$ is an isolated fixed point or not. In the latter case, the zeroes of $X$ form a null geodesic, and the resulting metric is a kind of plane wave, whose conformal vector fields can be determined. Interesting studies of the link between homothetic and conformal vector fields (and their fixed points) and the geometry the metrics can be found in [2, 4, 17, 18]. The above Theorem 3.1 and the special cases described in Theorem 3.6 and in Section 4, allow us to discuss the existence of homothetic fixed points for all three-dimensional Walker metrics, and gives a unified treatment for a large class of three-dimensional manifolds, where all different behaviours can occur, from metrics with no proper homothetic vector fields, to cases where homothetic fixed points occur and can be explicitly determined.

We now turn our attention to curvature collineations and prove the following.
Theorem 3.4. Let $X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y$ be an arbitrary smooth vector field on the strictly Walker manifold $(M, g_f)$, where $g_f$ is described as in (2.1). Then:

(i) $X$ is a Ricci collineation if and only if one of the following cases occurs:

(a) $f$ is arbitrary and

$$X_2 = -\frac{2f_1(y)\partial_{xx}^2 f + f_1(y)\partial_{xyy}^3 f}{\partial_{xxx}^3 f}, \quad X_3 = f_1(y),$$

where $f_1$ is an arbitrary smooth function on $M$, and the Ricci collineation is defined in the open subset where $\partial_{xxx}^3 f \neq 0$.

(b) $f = f_1(y)x^2 + f_2(y)x + f_3(y)$ and

$$X_3 = \frac{c_1}{\sqrt{|f_1(y)|}}.$$

(ii) $X$ is a curvature collineation if and only if $X$ is a special Ricci collineation of one of the following types:

(a') type (a) with

- either $X_1 = X_1(y), X_2 = X_3 = 0$, or
- $f(x, y) = f_2(x)f_3(y) + f_4(y)x + f_5(y), f_1(y) = \frac{c_1}{\sqrt{|f_3(y)|}}$, and

$$X_1 = \frac{c_1f_1(y)}{2}f_3(y)\sqrt{|f_3(y)|}t + f_6(y).$$

(b') type (b) with $X_1 = 2f_4(y)t + \frac{c_1f_1(y)}{2}f_3(y)\sqrt{|f_3(y)|}t - \frac{1}{2}f_1'(y)x^2 - f_5'(y)x + f_6(y)$ and

$$X_2 = f_4(y)x + f_5(y).$$

Proof. Because of equation (2.5), a smooth vector field $X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y$ on a strictly Walker manifold $(M, g_f)$ is a Ricci collineation if and only if

$$\partial_{xx}^2 f \partial_t X_3 = 0, \quad \partial_{xx}^2 f \partial_x X_3 = 0, \quad 2\partial_{xx}^2 f \partial_y X_3 + \partial_{xxx}^3 f X_2 + \partial_{xyy}^3 f X_3 = 0.$$  (3.9)

As we already mentioned, we are always assuming that $\partial_{xx}^2 f \neq 0$. Consequently, from the first two equations in (3.9) we have $X_3 = X_3(y)$, and the third equation becomes

$$2\partial_{xx}^2 f X_3'(y) + \partial_{xxx}^3 f X_2 + \partial_{xyy}^3 f X_3(y) = 0.$$  (3.10)

In the open subset where $\partial_{xxx}^3 f \neq 0$, from the above equation (3.10) we get at once the case (a). Case (b) is obtained as a special solution of (3.10), assuming that $\partial_{xxx}^3 f = 0$.

We then consider curvature collineations, starting from an arbitrary Ricci collineation as described in cases (a) and (b) and requiring the additional condition $\mathcal{L}_X R = 0$. Calculations are of the same kind for all these cases. For this reason, we report the details only for case (b).

So, consider a Ricci collineation $X = X_1 \partial_t + X_2 \partial_x + \frac{c_1}{\sqrt{|f_3(y)|}} \partial_y$, where $f = f_1(y)x^2 + f_2(y)x + f_3(y)$. In particular, calculating the condition $\mathcal{L}_X R = 0$ on the pairs of coordinate vector fields $\partial_t$, $\partial_x$, $\partial_y$, we find that $X$ is a curvature collineation if and only if the following equations hold:

$$\partial_y X_2 + \partial_x X_1 = \partial_t X_2 = 0, \quad \frac{c_1f_1'(y)}{2\sqrt{|f_1(y)|}} + f_1(y)(2\partial_x X_2 - \partial_t X_1) = 0.$$  (3.11)

It easily follows from the first of the above equations that $\partial_x X_1$ and $X_2$ are functions of the variables $(x, y)$. Since $f_1(y) \neq 0$, differentiating with respect to $x$ the second of the above equations we get $\partial_{xx}^2 X_2 = 0$ and so, $X_2 = f_4(y)x + f_5(y)$. Now, again the second equation gives $X_1 = \frac{c_1f_1'(y)}{2f_4(y)\sqrt{|f_3(y)|}} + f_6(y)t + f_6(x, y)$. Then, since $\partial_x X_1 = -\partial_y X_2$, we conclude that $f_6(x, y) = -\frac{1}{2}f_1'(y)x^2 - f_5'(y)x + f_7(y)$ and this ends the proof. □
Observe that taking $X_2 = X_3 = 0$, all equations in (3.9) are satisfied. Therefore, $X = X_1 \partial_t$ is a Ricci collineation for any arbitrary smooth function $X_1 = X_1(t, x, y)$, and (by case (a)') a curvature collineation for any smooth function $X_1 = X_1(y)$. This implies at once the following.

**Corollary 3.5.** For any strictly Walker three-manifold $(M, g_f)$, the Lie algebras of smooth Ricci collineations and smooth curvature collineations are infinite-dimensional. In particular, each of these spaces admits proper Ricci and curvature collineations.

We end this section calculating the symmetries of a locally conformally flat strictly Walker three-manifold. By direct calculations of the Cotton tensor of a strictly Walker three-manifold $(M, g_f)$ (see also [12]), it is easily seen that this manifold is locally conformally flat if and only if $\partial^3_{xxy} f$ vanishes identically, that is, when the defining function is of the form $f(x, y) = p(y)x^2 + q(y)x + r(y)$ (with $p(y) \neq 0$ in order to avoid the flat case). We now prove the following.

**Theorem 3.6.** Let $X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y$ be an arbitrary smooth vector field on a conformally flat strictly Walker manifold $(M, g_f)$, where $g_f$ is described as in (2.1) with $f(x, y) = p(y)x^2 + q(y)x + r(y)$ ($p(y) \neq 0$). Then, $X$ is:

i) a Killing vector field if and only if $(c_1 y + c_2)p(y)' + 2c_1 p(y) = 0$ and $X$ is described as in (3.1), with $f_1$, $f_2$ explicitly determined as solutions of

\[
2f_1(y)'' - 2c_1 q(y) - (c_1 y + c_2)q(y)' - 2f_1(y)p(y) = 0,
\]
\[
2f_2(y)' + 2c_1 r(y) + (c_1 y + c_2)r(y)' + f_1(y)q(y) = 0.
\]

ii) a homothetic, non-Killing vector field if and only if $(c_1 y + c_2)p(y)' + 2c_1 p(y) = 0$ and $X$ is described as in (3.3), with $f_1$, $f_2$ explicitly determined as solutions of

\[
2f_1''(y) + (\frac{\eta}{2} - 2c_1)q(y) - (c_1 y + c_2)q(y)' - 2f_1(y)p(y) = 0,
\]
\[
2f_2''(y) + (2c_1 - \eta)r(y) + (c_1 y + c_2)r(y)' + f_1(y)q(y) = 0, \quad \eta \neq 0.
\]

iii) a proper affine Killing vector field if and only if $(c_2 y + c_3)p(y)' + 2c_2 p(y) = 0$ and $X$ is described as in (3.5), with $f_1$, $f_2$ explicitly determined as solutions of

\[
2f_1''(y) + \frac{c_1 - 3c_2}{2} q(y) - (c_2 y + c_3)q(y)' - 2f_1(y)p(y) = 0,
\]
\[
2f_2''(y) + (c_2 - c_1)r(y) + (c_2 y + c_3)r(y)' + f_1(y)q(y) + c_4 = 0.
\]

iv) a Ricci collineation if and only if $X_3 = \frac{c_1}{\sqrt{|p(y)|}}$, where $c_1$ is a real constant.

v) a curvature collineation if and only if

\[
X_1 = 2f_1(y)t + \frac{c_1 p'(y)}{2p(y)\sqrt{|p(y)|}} - \frac{1}{2} f_1'(y)x^2 - f_2'(y)x + f_3(y),
\]
\[
X_2 = f_1(y)x + f_2(y), \quad X_3 = \frac{c_1}{\sqrt{|p(y)|}},
\]

where $f_1(y)$ and $f_2(y)$ are arbitrary smooth functions on $M$.

**Proof.** Let $(M, g_f)$ be a conformally flat strictly Walker manifold of dimension three, where $g_f$ is described by the relation (2.1). As explained above, the function $f$ satisfies $f(x, y) = p(y)x^2 + q(y)x + r(y)$ ($p(y) \neq 0$), where $p(y)$, $q(y)$ and $r(y)$ are arbitrary smooth functions on $M$. We then choose a smooth vector field $X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y$, where $X_1$, $X_2$ and $X_3$ are arbitrary
smooth functions on $M$. By Theorem 3.1, $X$ is a homothetic vector field if and only if satisfies equations (3.3) and (3.4). Equation (3.4) for the above function $f(x,y)$ gives

$$
(2c_1 p(y) + (c_1 y + c_2)p(y)^{\prime}) x^2
- \left(2f_1(y)^{\prime\prime} + \left(\frac{\eta}{2} - 2c_1\right) q(y) - (c_1 y + c_2)q(y)^{\prime} - 2f_1(y)p(y)\right) x
+ (2c_1 - \eta)r(y) + 2f_2(y)^{\prime} + (c_1 y + c_2)r(y)^{\prime} + f_1(y)q(y) = 0.
$$

This equation immediately proves the second statement, since the coefficients of $x$ and its powers must vanish, in order to satisfy it identically. The first statement now follows by setting $\eta = 0$ in the equations of homothetic vector fields.

With regard to affine Killing vector fields, $X$ must satisfy equations (3.5) and (3.6). So by straightforward calculations, the functions $p(y)$, $q(y)$ and $r(y)$ must satisfy

$$
(2c_2 p(y) + (c_2 y + c_3)p(y)^{\prime}) x^2
- \left(2f_1^{\prime\prime}(y) - (c_2 y + c_3)q'(y) - 2f_1(y)p(y) + \frac{c_1 - 3c_2}{2}q(y)\right) x
+ (2f_2(y) + (c_2 y + c_3)r(y) + f_1(y)q(y) + (c_2 - c_1)r(y) + c_4) = 0,
$$

which leads to the third statement. Assertions (iv) and (v) are direct consequences of the cases (b) and (b)' of Theorem 3.4, respectively.

\[\Box\]

4 Symmetries of homogeneous Lorentzian three-manifolds with recurrent curvature

We reported the classification of homogeneous Lorentzian three-manifolds with recurrent curvature in Theorem 2.1. We shall now completely describe the symmetries of these manifolds. We start with the following.

Theorem 4.1. Let $(M,g_f)$ be a homogeneous three-dimensional Lorentzian strictly Walker manifold of type $N_0$, that is, determined by $f(x,y) = \frac{-2e^{bx}}{b^2}$, $b \neq 0$.

An arbitrary smooth vector field $X = X_1\partial_t + X_2\partial_x + X_3\partial_y$ on $M$:

- is homothetic (equivalently, Killing) if and only if
  $$
  X_1 = c_1 t + c_2, \quad X_2 = \frac{2c_1}{b}, \quad X_3 = -c_1 y + c_3.
  $$

- is affine if and only if
  $$
  X_1 = c_1 t + c_2 + c_4 y, \quad X_2 = \frac{2c_1}{b}, \quad X_3 = -c_1 y + c_3.
  $$

- is a Ricci collineation if and only if $X_1$ is arbitrary and
  $$
  X_2 = -\frac{2}{b}f_1'(y), \quad X_3 = f_1(y).
  $$

- is a curvature collineation if and only if
  $$
  X_1 = f_2(y) - f_1'(y)t + \frac{2}{b}f_1''(y)x, \quad X_2 = -\frac{2}{b}f_1'(y), \quad X_3 = f_1(y).
  $$

In the above equations, $c_j$ are real constants and $f_1(y)$ an arbitrary smooth function.
Theorem 4.2. Let \((M, g_f)\) be a homogeneous three-dimensional Lorentzian strictly Walker manifold of type \(\mathcal{P}_c\), that is, determined by \(f(x, y) = -x^2 \alpha(y)\), where
\[
\alpha'(y) = c \alpha(y)^{\frac{3}{2}}, \quad \alpha(y) > 0.
\]
Let \(h(y)\) denote a smooth function explicitly determined from \(\alpha(y)\) by equation
\[
h''(y) + \alpha(y)h(y) = 0.
\]
An arbitrary smooth vector field \(X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y\) on \(M\):
- is Killing if and only if
\[
X_1 = -h'(y)x + c_1, \quad X_2 = h(y), \quad X_3 = 0.
\]
• is homothetic if and only if
\[ X_1 = -h'(y)x + c_1 + \eta t, \quad X_2 = h(y) + \frac{\eta}{2} x, \quad X_3 = 0. \]
• is affine if and only if
\[ X_1 = -h'(y)x + c_1 + c_2 t + c_3 y, \quad X_2 = h(y) + \frac{c_2}{2} x, \quad X_3 = 0. \]
• is a Ricci collineation if and only if \( X_1, X_2 \) are arbitrary and \( X_3 = \frac{c_1}{\sqrt{|\alpha(y)|}} \).
• \( X \) is a curvature collineation if and only if
\[ X_1 = -\frac{1}{2} f_1'(y)x^2 - f_2'(y)x + \left(2 f_1(y) + \frac{c_1 c}{2}\right)t + f_3(y), \]
\[ X_2 = f_1(y)x + f_2(y), \quad X_3 = \frac{c_1}{\sqrt{|\alpha(y)|}}. \]

In the above equations, \( f_i(y) \) are arbitrary smooth functions and \( c_j \) are real constants.

**Theorem 4.3.** Let \((M, g_f)\) be a homogeneous three-dimensional Lorentzian strictly Walker manifold of type \( CW_\varepsilon \), that is, determined by \( f(x, y) = -\varepsilon x^2 \). Consider the functions \( s(y), t(y) \) given by
\[ s(y) = \begin{cases} 
    c_2 \sin(y) - c_1 \cos(y) & \text{if } \varepsilon = 1, \\
    c_1 e^{-y} - c_2 e^y & \text{if } \varepsilon = -1,
\end{cases} \quad t(y) = \begin{cases} 
    c_1 \sin(y) + c_2 \cos(y) & \text{if } \varepsilon = 1, \\
    c_1 e^{-y} + c_2 e^y & \text{if } \varepsilon = -1.
\end{cases} \]

Then, an arbitrary smooth vector field \( X = X_1 \partial_t + X_2 \partial_x + X_3 \partial_y \) on \( M \):
• is Killing if and only if
\[ X_1 = s(y)x + c_3, \quad X_2 = t(y), \quad X_3 = c_4, \]
• is homothetic if and only if
\[ X_1 = s(y)x + c_3 + \eta t, \quad X_2 = t(y) + \frac{\eta}{2} x, \quad X_3 = c_4, \]
• is affine if and only if
\[ X_1 = s(y)x + c_3 + c_5 t + c_6 y, \quad X_2 = t(y) + \frac{c_5}{2} x, \quad X_3 = c_4, \]
• is a Ricci collineation if and only if \( X_3 = c_1 \).
• is a curvature collineation if and only if
\[ X_1 = 2f_1(y)t - \frac{1}{2} f_1'(y)x^2 - f_2'(y)x + f_3(y), \quad X_2 = f_1(y)x + f_2(y), \quad X_3 = c_1. \]

In the above equations, \( f_i(y) \) are arbitrary smooth functions on \( M \) and \( c_j \) real constants.

It is well known that for three-dimensional manifolds and locally conformally flat manifolds, the curvature is completely determined by its Ricci curvature. However, as a consequence of Theorem 3.6 (in particular, Theorems 4.2 and 4.3), we find explicit examples of three-dimensional locally conformally flat spaces, for which Ricci and curvature collineations are not equivalent.
Acknowledgements

First author partially supported by funds of the University of Salento and MIUR (PRIN). Second author partially supported by funds of the University of Payame Noor. The authors wish to thank the anonymous referees for their valuable suggestions and comments.

References

[1] Aichelburg P.C., Curvature collineations for gravitational pp waves, *J. Math. Phys.* 11 (1970), 2458–2462.
[2] Alekseevski D., Self-similar Lorentzian manifolds, *Ann. Global Anal. Geom.* 3 (1985), 59–84.
[3] Batat W., Calvaruso G., De Leo B., Homogeneous Lorentzian 3-manifolds with a parallel null vector field, *Balkan J. Geom. Appl.* 14 (2009), 11–20.
[4] Beem J.K., Proper homothetic maps and fixed points, *Lett. Math. Phys.* 2 (1978), 317–320.
[5] Brozos-Vázquez M., García-Río E., Gilkey P., Nikčević S., Vázquez-Lorenzo R., The geometry of Walker manifolds, *Synthesis Lectures on Mathematics and Statistics*, Vol. 5, Morgan & Claypool Publishers, Williston, VT, 2009.
[6] Calvaruso G., Zaeim A., Invariant symmetries on non-reductive homogeneous pseudo-Riemannian four-manifolds, *Rev. Mat. Complut.* 28 (2015), 599–622.
[7] Calvaruso G., Zaeim A., Geometric structures over four-dimensional generalized symmetric spaces, *Collect. Math.*, to appear.
[8] Calvino-Louzao E., Seoane-Bascoy J., Vázquez-Abal M.E., Vázquez-Lorenzo R., Invariant Ricci collineations on three-dimensional Lie groups, *J. Geom. Phys.* 96 (2015), 59–71.
[9] Camci U., Hussain I., Kucukakca Y., Curvature and Weyl collineations of Bianchi type V spacetimes, *J. Geom. Phys.* 59 (2009), 1476–1484.
[10] Camci U., Sharif M., Matter collineations of spacetime homogeneous Gödel-type metrics, *Classical Quantum Gravity* 20 (2003), 2169–2179, gr-qc/0306129.
[11] Carot J., da Costa J., Vaz E.G.L.R., Matter collineations: the inverse “symmetry inheritance” problem, *J. Math. Phys.* 35 (1994), 4832–4838.
[12] Chaichi M., García-Río E., Vázquez-Abal M.E., Three-dimensional Lorentz manifolds admitting a parallel null vector field, *J. Phys. A: Math. Gen.* 38 (2005), 841–850.
[13] Flores J.L., Parra Y., Percoco U., On the general structure of Ricci collineations for type B warped spacetimes, *J. Math. Phys.* 45 (2004), 3546–3557, gr-qc/0405133.
[14] García-Río E., Gilkey P.B., Nikčević S., Homogeneity of Lorentzian three-manifolds with recurrent curvature, *Math. Nachr.* 287 (2014), 32–47, arXiv:1210.7764.
[15] Hall G., Symmetries of the curvature, Weyl conformal and Weyl projective tensors on 4-dimensional Lorentz manifolds, in Proceedings of the International Conference “Differential Geometry – Dynamical Systems” (DGDS-2007), BSG Proc., Vol. 15, Geom. Balkan Press, Bucharest, 2008, 89–98.
[16] Hall G.S., Symmetries and curvature structure in general relativity, *World Scientific Lecture Notes in Physics*, Vol. 46, World Scientific Publishing Co., Inc., River Edge, NJ, 2004.
[17] Hall G.S., Capocci M.S., Classification and conformal symmetry in three-dimensional space-times, *J. Math. Phys.* 40 (1999), 1466–1478.
[18] Hall G.S., Low D.J., Pulham J.R., Affine collineations in general relativity and their fixed point structure, *J. Math. Phys.* 35 (1994), 5930–5944.
[19] Hall G.S., Roy I., Vaz E.G.L.R., Ricci and matter collineations in space-time, *Gen. Relativity Gravitation* 28 (1996), 299–310.
[20] Kühnel W., Rademacher H.-B., Conformal Ricci collineations of space-times, *Gen. Relativity Gravitation* 33 (2001), 1905–1914.
[21] Levichev A.V., Methods for studying the causal structure of homogeneous Lorentz manifolds, *Sib. Math. J.* 31 (1990), 395–408.
[22] Tsamparlis M., Apostolopoulos P.S., Ricci and matter collineations of locally rotationally symmetric spacetimes, *Gen. Relativity Gravitation* 36 (2004), 47–69, gr-qc/0309034.
[23] Walker A.G., On parallel fields of partially null vector spaces, *Quart. J. Math. Oxford Ser.* 20 (1949), 135–145.