COMPARISON OF DIFFERENT GOODNESS-OF-FIT TESTS

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Abstract

Various distribution free goodness-of-fit test procedures have been extracted from literature. We present two new binning free tests, the univariate three-region-test and the multivariate energy test. The power of the selected tests with respect to different slowly varying distortions of experimental distributions are investigated. None of the tests is optimum for all distortions. The energy test has high power in many applications and is superior to the $\chi^2$ test.

1 INTRODUCTION

Goodness-of-fit (gof) tests are designed to measure the compatibility of a random sample with a theoretical probability distribution function (pdf). The null hypothesis $H_0$ is that the sample follows the pdf. Under the assumption that $H_0$ applies, the fraction of wrongly rejected experiments - the probability of committing an error of the first kind - is fixed to typically a few percent. A test is considered powerful if the probability of accepting $H_0$ when $H_0$ is wrong - the probability of committing an error of the second kind - is low. Of course, without specifying the alternatives, the power cannot be quantified.

A discrepancy between a data sample and the theoretical description can be of different origin. The problem may be in the theory which is wrong or the sample may be biased by measurement errors or by background contamination. In natural sciences we mainly have the latter situation. Even though the statistical description is the same in both cases the choice of the
specific test may be different. In our applications we are mainly confronted with “slowly varying” deviations between data and theoretical description whereas in other fields where for example time series are investigated, “high frequency” distortions are more likely.

Goodness-of-fit tests are based on classical statistical methods and are closely related to classical interval estimation, but they contain also Bayesian elements. Those, however, are only related with some prejudice on the alternative hypothesis which affects the purity of the accepted decisions and not the error of the first kind.

The power of one dimensional tests is not always invariant against transformations of the variates. In more than one dimension (number of variates), an invariant description is not possible.

Tests are classified in distribution dependent and distribution free tests. The former are adapted to special pdfs like Gaussian, exponential or uniform distributions. We will restrict our discussion mainly to distribution free tests and tests which can be adapted to arbitrary distributions. Here we distinguish tests applied to binned data and binning free tests. The latter are in principle preferable but so far they are almost exclusively limited to one dimensional distributions. A further distinction concerns the alternative hypothesis. Usually, it is not restricted but there exist also tests where it is parametrized.

Physicists tend to be content with $\chi^2$ tests which are not necessarily optimum in all cases. A very useful and comprehensive survey of goodness-of-fit tests can be found in Ref. [1] from 1986. Since then, some new developments have occurred and the increase in computing power has opened the possibility to apply more elaborate tests.

In Section 2 we summarize the most important tests. To keep this article short we do not discuss tests based on the order statistic and spacing tests. In Section 3 we introduce two new tests, the three region test and the energy test. To compare the tests we apply them in Section 4 to some specific alternative hypotheses. We do not consider explicitly composite hypotheses.

2 SOME RELEVANT TESTS
2.1 Chi-squared Test

The $\chi^2$ test is closely connected to least square fits with the difference that the hypothesis is fixed. The test statistic is

$$\chi^2 = \sum_{i=1}^{B} \frac{(Y_i - t_i)^2}{\delta_i^2}$$

with the random variable $Y_i$, $t_i$ the expectation $E(Y_i)$ and $\delta_i^2$ the expectation $E((Y_i - t_i)^2)$. Obviously the expectation value of $\chi^2$ is

$$E(\chi^2) = B$$

In the Gaussian approximation $Y_i$ follows a Gaussian with mean $t_i$ and variance $\delta_i^2$ and the test statistic follows a $\chi^2$ distribution function $F_B(\chi^2)$ with $B$ degrees of freedom. The probability of an error of the first kind $\alpha$ (significance level, p-value) defines $\chi^2_0$ with $F_B(\chi^2_0) = 1 - \alpha$. The null hypothesis is rejected if in an actual experiment we find $\chi^2 > \chi^2_0$.

We obtain the Pearson test when the random variables $N_i$ are Poisson distributed.

$$\chi^2 = \sum_{i=1}^{B} \frac{(N_i - t_i)^2}{t_i}$$

In the large sample limit the test statistic $\chi^2$ has approximately a $\chi^2$ distribution with $B$ degrees of freedom. When we have a histogram, where a total of $N$ events are distributed according to a multinomial distribution among the $B$ bins with probabilities $p_i$ we get

$$\chi^2 = \sum_{i=1}^{B} \frac{(N_i - Np_i)^2}{Np_i}$$

which again follows asymptotically a $\chi^2$ distribution, this time with $B - 1$ degrees of freedom. The reduced number of degrees of freedom is due to the constraint $\Sigma N_i = N$.

Nowadays, the distribution function of the test statistic can be computed numerically without much effort. The $\chi^2$ test then can also be applied to small samples. The Gaussian approximation is no longer required.

The $\chi^2$ test is very simple and needs only limited computational power. A big advantage compared to most of the other methods is that is can be applied to multidimensional histograms. There are however also serious drawbacks:
• Its power in detecting slowly varying deviations of a histogram from predictions is rather poor due to the neglect of possible correlations between adjacent bins.

• Binning is required and the choice of the binning is arbitrary.

• When the statistics is low or the number of dimensions is high, the event numbers per bin may be low. Then the asymptotic properties are no longer valid and systematic deviations are hidden by statistical fluctuations.

There are proposals to fix the bin widths by the requirement of equal number of expected entries per bin. This is not necessarily the optimum choice \[2\]. Often there are outliers in regions where no events are expected which would be hidden in wide bins.

For the number of bins a dependence on the sample size \(n\)

\[ B = 2n^{2/5} \]

is proposed in Ref. \[3\]. Our experience is that in most experiments the number of bins is chosen too high. The sensitivity to slowly varying deviations roughly goes with \(B^{-1/4}\) \[2\]. In multidimensional cases the power of the test often can be increased by applying it to the marginal distributions.

There is a whole class of \(\chi^2\) like tests. Many studies can be found in the literature. The reader is referred to Ref. \[3\].

### 2.2 Binning-free empirical distribution function tests

The tests described in this section have been taken from the article by Stephens in Ref. \[1\].

Supposing that a random sample of size \(n\) is given, we form the order statistic \(X_1 < X_2 < \ldots < X_n\). We consider the empirical distribution function (EDF)

\[ F_n(x) = \frac{\# \text{ of observations } \leq x}{n} \]

or

\[
\begin{align*}
F_n(x) &= 0 \quad x < X_1 \\
F_n(x) &= \frac{i}{n} \quad X_i \leq x < X_{i+1} \\
F_n(x) &= 1 \quad X_n \leq x
\end{align*}
\]
Figure 1: Comparison of empirical and theoretical distributions

$F_n(x)$ is a step function which is to be compared to the distribution $F(x)$ corresponding to $H_0$.

The EDF is consistent and unbiased. The tests discussed in this section are invariant under transformation of the random variable. Because of this feature, we can transform the distribution to the uniform distribution and restrict our discussion to the latter.

2.2.1 Probability integral transformation

The probability integral transformation (PIT)

$$Z = F(X)$$

transforms a general pdf $f(X)$ of $X$ into the uniform distribution $f^*(Z)$ of $Z$.

$$f^*(Z) = 1; \quad 0 \leq Z \leq 1$$

$$F^*(Z) = Z$$

The underlying idea of this transformation is that the new EDF of $Z$, $F^*(Z)$ is extremely simple and that it conserves the distribution of the test quantities discussed in this section. It is easily seen that
\[ F_n(x) - F(x) = F^*_n(z) - z \]

Note, however, that the PIT does not necessarily conserve all interesting features of the gof problem. Resolution effects are washed out and for example in a lifetime distribution, an excess of events at small and large lifetimes may be judged differently but are treated similarly after a PIT. It is not logical to select specific gofs for specific applications but to transform all kinds of pdfs to the same uniform distribution. The PIT is very useful because it permits standardization but one has to be aware of its limitations.

2.2.2 Supremum statistics

The maximum positive (negative) deviation of \( F_n(x) \) from \( F(x) \) \( D_+ \) (\( D_- \)) (see Fig. 1) are used as tests statistics. Kolmogorov (Kolmogorov-Smirnov test) has proposed to use the maximum absolute difference. Kuiper uses the sum \( V = D_+ + D_- \). This test statistic is useful for observations “on the circle” for example for azimuthal distributions where the zero angle is a matter of definition.

\[
\begin{align*}
D_+ &= \sup_x \{ F_n(x) - F(x) \} \\
D_- &= \sup_x \{ F(x) - F_n(x) \} \\
D &= \sup_x \{|F_n(x) - F(x)|\} \quad \text{(Kolmogorov)} \\
V &= D_+ + D_- \quad \text{(Kuiper)}
\end{align*}
\]

The supremum statistics are invariant under the PIT.

2.2.3 Integrated deviations - quadratic statistics

The Cramer-von Mises family of tests measures the integrated quadratic deviation of \( F_n(x) \) from \( F(x) \) suitably weighted by a weighting function \( \psi \):

\[
Q = n \int_{-\infty}^{\infty} [F(x) - F_n(x)]^2 \psi(x) dF
\]

In the standardized form we have

\[
Q = n \int_0^1 |z - F^*_n(z)|^2 \psi(z) dz \quad \text{(1)}
\]
Since the construction of $F_n^*(Z)$ includes already an integration, $F_n^*(z_i)$ and $F_n^*(z_k)$ are not independent and the additional integration in Equation \( \square \) is not obvious.

With $\psi_{CvM} = 1$ we get the Cramer-von Mises statistic $W^2$ and $\psi_{AD} = [z(1-z)]^{-1}$ leads to the Anderson-Darling statistic $A^2$.

$$W^2 = n \int_0^1 [z - F_n^*(z)]^2 dz \quad \text{(Cramer-von Mises)}$$

$$A^2 = n \int_0^1 \frac{[z - F_n^*(z)]^2}{z(1-z)} dz \quad \text{(Anderson - Darling)}$$

The Anderson-Darling statistic $A^2$ weights strongly deviations near $z = 0$ and $z = 1$. This is justified because there the experimental deviations are small due to the constraints $[z - F_n^*(z)] = 0$ at $z = 0$ and $z = 1$.

Watson has proposed a quadratic statistic on the circle:

$$U^2 = n \int_0^1 \left\{ F_n^*(z) - z - \int_0^1 [F_n^*(z) - z] \, dz \right\}^2 dz \quad \text{(Watson)}$$

### 2.3 The Neyman statistic test

This test is different from all previously discussed tests. It parametrizes the alternative hypothesis and applies the likelihood ratio test. The alternative hypothesis corresponds to a pdf of the exponential family:

$$g_k(z) = C(\theta_1, \theta_2...\theta_k) \exp \left[ \sum_{i=1}^k \theta_i \pi_i(z) \right]$$

$g_k(z)$ are smooth alternatives to uniformity. The functions $\pi_i$ are Legendre polynomials of order $i$, $\theta_i$ are free parameters and $C$ is a normalization function. The number $k$ of parameters is selected by the user.

The likelihood ratio leads to the test statistic

$$N_k = \frac{1}{n} \sum_{i=1}^k \left( \sum_{j=1}^n \pi_i(z_j) \right)^2$$

Asymptotically, for large values of $N_k$, $N_k$ is distributed according to the $\chi^2$ distribution with $k$ degrees of freedom.

7
3 NEW TESTS

3.1 Three region test

Often experimental distributions are biased by an excess or lack of events in certain regions of the random variable. We have designed a test which subdivides the variable space into three pieces, containing \( n_1, n_2, n_3 = n - n_1 - n_2 \) events, such that the deviation between data and prediction from \( H_0 \) is maximum. The test quantity is

\[
O = \sup_{n_1, n_2} \left\{ w_1(n_1 - np_1)^2 + w_2(n_2 - np_2)^2 + w_3(n_3 - np_3)^2 \right\}
\]

where \( np_k \) are the expectation values and \( w_k \) weights depending on \( np_k \). The specific choice

\[
w_k = \frac{1}{np_k}
\]

\[
O_\chi = \sup_{n_1, n_2} \left\{ \frac{(n_1 - np_1)^2}{np_1} + \frac{(n_2 - np_2)^2}{np_2} + \frac{(n_3 - np_3)^2}{np_3} \right\}
\]

maximizes \( \chi^2 \) of the three bins. In the comparison below we have chosen weights equal to one.

Of course the test can be generalized to a higher number of subregions.

3.2 Minimum energy test

3.2.1 The idea

Let us assume that we have a continuous charge distribution \( \rho(\vec{r}) \) of positive electric charges and a sample of negative point charges with total charge equal to minus the integrated positive charge. The potential energy is minimum when the negative point charges follow \( \rho \). Then, up to effects due to the discrete nature of the point charges, the charge density is zero everywhere. Any displacement of charges would increase the energy. We use this property to construct a binning free test procedure.

We simulate the theoretical distribution by \( m \) charges of charge \( 1/m \) each. Usually, these charges are distributed using a Monte Carlo simulation. To the \( n \) experimental sample points (data points) we associate charges \(-1/n\). The test quantity \( \phi \) corresponds to the potential energy. It contains two terms
\(\phi_1, \phi_2\) corresponding to the interaction of the experimental charges with each other and to the interaction of the experimental charges with the positive simulation charges.

\[
\phi_1 = \frac{1}{n^2} \sum_{i<j} R(d_{ij}) \quad (2)
\]

\[
\phi_2 = -\frac{1}{nm} \sum_{i,j} R(t_{ij}) \quad (3)
\]

\[
\phi = \phi_1 + \phi_2 \quad (4)
\]

Here \(d_{ij}\) is the distance between two data points and \(t_{ij}\) is the distance between a data point and a simulation point and \(R\) is a correlation function defined below. The sums run over all combinations.

Remark: The minimum energy requirement for the equality of experimental and theoretical distribution is strictly correct only when the number \(m\) of simulation charges is equal to the number \(n\) of experimental charges. For the general case with a continuous theoretical distribution or simulation sample and experimental sample of different size, the optimum agreement of the two distributions is not well defined and there is a slight dependence of the minimum energy configuration on the correlation function. This is however a purely academic problem, the test statistic \(\phi\) remains a powerful indicator for an incompatibility of the experimental sample with \(H_0\).

### 3.2.2 The correlation function

We note that the minimum energy configuration does not depend on the application of the one-over-distance power law of electrostatics. We may apply a wide class of correlation functions \(R(r)\) with the only requirement that \(R\) has to decrease monotonically with the distance \(r\).

We have investigated three different types of correlation functions, power laws, a logarithmic dependence and Gaussians.

\[
R_\kappa(r) = \frac{1}{r^\kappa} \quad (5)
\]

\[
R_l(r) = -\ln r \quad (6)
\]

\[
R_s(r) = e^{-r^2/(2\sigma^2)} \quad (7)
\]

The first type is motivated by the analogy to electrostatics, the second is long range and the third emphasizes a limited range for the correlation
between different points. The power \( \kappa \) of the denominator in Equ. 5 and the parameter \( s \) in Equ. 6 may be chosen differently for different dimensions of the sample space and different applications. For long range distortions a small value of \( \kappa \) around 0.1 is recommended. For short range deviations the test quantity with larger values around 0.3 is more sensitive.

The inverse power law and the logarithm have a singularity at \( r \) equal to zero. Very small distances, however, should not be weighted too strongly since distortions with sharp peaks are not expected and usually inhibited by the finite experimental resolution. We eliminate the singularity by introducing a lower cutoff \( d_{\text{min}} \) for the distances \( d \) and \( t \). Distances less than \( d_{\text{min}} \) are replaced by \( d_{\text{min}} \). The value of this parameter is not critical, it should be of the order of the average distance \( d \) in the regions where the \( f_0 \) is maximum and not less than the experimental resolution.

The energy test with Gaussian correlation function is closely related to the Pearson \( \chi^2 \) test. A more detailed description of the energy test is presented in Ref. [4].

### 3.3 Comparison of uni-variate tests

We have tested the null hypothesis of a uniform distribution in the interval \([0, 1]\) using a uniform distribution contaminated by the background distributions displayed in Figure 2.

Background hypothesis A modifies the mean, hypotheses B, C change the variance of \( H_0 \).

The power of various tests described above is presented in Figure 3.

As expected, none of the tests is optimum for all kind of distortions. Several tests perform better than the \( \chi^2 \) test. The Neyman test, the Anderson-Darling test and the Kolmogorov-Smirnov test are sensitive to a shift in the mean. Anderson’s test detects especially deviations at the borders of the interval. Watson’s and Kuiper’s tests are useful for the detection of distortions of the variance. The two new tests compare favorably with the standard ones.

### 4 MULTIVARIATE TESTS

The Mardia test [5] and the Neyman smooth test [6] can be used to investigate two-dimensional Gaussian distributions. The only distribution free test
Figure 2: Different types of background distributions

Figure 3: Histograms show the powers of different tests at 5% level for the sample size n=100.
known to us which is independent of the dimensions of the variate space is the $\chi^2$ test. A generalized Kolmogorov-Smirnov test depends on the ordering scheme of the variates. The binning free energy test developed by us is also independent of the number of variates, however the distribution of test statistic has to be computed for the specific sample distribution under study.

### 4.1 Comparison of multivariate tests

We have used a two-dimensional Gaussian null hypothesis and contaminated the sample with the background distributions shown in Figure 4.

The power of the two Mardia tests, the Neyman smooth test and the energy test with logarithmic and Gaussian correlation function is presented in Figure 5.

In most cases the two energy tests perform better than the alternatives even though those have been designed for a Gaussian null hypothesis.

### 4.2 Example: Comparing experimental data to a Monte Carlo prediction

In Figure 6, left hand side, we compare the position and momentum of a few $J/\psi$ decay tracks to a Monte Carlo simulation. The right hand plot compares the energy computed from the distribution on the left hand side to a Monte Carlo simulation of the null hypothesis. The experimental point, indicated by the arrow, is larger than all Monte Carlo values. Apparently, the data do not follow the prediction.
Figure 5: Power comparisons of tests of bivariate normality for the sample size $n=200$ at 5% significance level.

Figure 6: Comparison of experimental distribution (squares) with Monte Carlo simulation (dots). The experimental energy computed from the scatter plot (left) is compared to a Monte Carlo simulation of the experiment (right).
5 CONCLUSIONS

The $\chi^2$ test suffers from the requirement to choose a binning. In one dimension it should be replaced by the well established binning free tests like the Kolmogorov test. The choice of a specific test has to depend on the expected kind of possible distortion of the theoretical distribution. For a localized background we advise to use the new three region test. For multivariate applications the new energy test is an attractive alternative to the $\chi^2$ test.

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