Jet properties from $\pi^\pm - h^\pm$ correlation in $p+p$ and $d+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV

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Abstract. We review recent results on the charged pion - charged hadron correlation in $p+p$ and $d+Au$ collisions as measured by the PHENIX Collaboration. Properties of di-jet system, such as the jet shape, associated hadron yield per trigger pion, and the underlying event are extracted statistically from the $\pi^\pm - h^\pm$ correlation function in $\Delta\phi$ and $\Delta\eta$. For jet triggered with high $p_T$ pions ($p_T > 5$ GeV/c), no apparent differences in the jet properties are seen between $p+p$ and $d+Au$.

1. Introduction

The technique of two particle correlation in relative azimuth ($\Delta\phi$) and pseudorapidity ($\Delta\eta$) is an useful tool to access the (di-)jet properties in heavy-ion collisions. Comparing with the traditional full jet reconstruction method, the two particle correlation method is relatively insensitive to the level of the underlying event, thus can probe soft jets ($\lesssim 5$ GeV/c); combining with event mixing technique, it can also be used for detectors with limited acceptance.

To leading order in QCD, high $p_T$ jets are produced back-to-back in azimuth. This back-to-back correlation, however, is smeared by the fragmentation process and the initial and final state radiation, which lead to a typical di-hadron correlation function in $\Delta\phi$ as shown schematically in Figure 1. The associated hadron yield per trigger $\pi^\pm$ (conditional yield or CY) can be parameterized by a constant plus a double gauss function,

$$
\frac{1}{N_0\text{trig}} \frac{dN}{d\Delta\phi} = B + \frac{N_S}{\sqrt{2\pi}\sigma_N} e^{-\frac{\Delta\phi^2}{2\sigma_N^2}} + \frac{N_A}{\sqrt{2\pi}\sigma_F} e^{-\frac{(\Delta\phi-\pi)^2}{2\sigma_F^2}},
$$

In this analysis, everything about the (di-)jet is extracted from this parameterization. The peaks in the same side ($\Delta\phi = 0$) and the away side ($\Delta\phi = \pi$) represent the intra-jet and di-jet correlation, respectively. The widths of the peaks are controlled by the jet fragmentation momentum $j_T$ and the parton transverse momentum $k_T$ [1, 2]: $\sigma_{\text{same}} \propto j_T$, $\sqrt{\sigma_{\text{away}}^2 - \sigma_{\text{same}}^2} \propto k_T$, where the subscript “$y$” represent the 1D projection in transverse plane; The integrals of the peaks, $N_S$ and $N_A$ give the total number of hadrons associated with the trigger hadrons in the same side and the away side; The pedestal level beneath the jet structure, $B$, represents contributions from the underlying event.

The presence of the medium can modify the di-hadron correlation. Recent results from RHIC [3, 4, 5] and SPS [6] indicate a complicated modification of the jet structure in both
Figure 1. Cartoon of the two particle $\Delta \phi$ correlation. The yield of hadron per trigger (Conditional Yield) has a di-jet part and a part corresponding to the underlying event.

the same side and the away side. The same side jet become elongated in $\Delta \eta$ but is relatively unmodified in $\Delta \phi$ direction. Meanwhile, the away side correlation become much broader and suppressed, which indicates strong interaction of the jets with the dense medium. In order to achieve quantitative understandings of the modifications of jets in the medium, one has to obtain an accurate baseline measurements of jet correlation in $p + p$ and $p + A$ collisions. Di-hadron correlation in $p + p$ collisions probes basic QCD effects such as jet fragmentation, initial and final state radiations; while correlation in $p + A$ collisions gives handle on the various initial state effects such as shadowing and jet broadening in cold nuclear medium.

We focus the physics discussions on three aspects of the $\pi^\pm - h^\pm$ correlations in Figure 1: jet shape, jet yield and the underlying event. But before doing that, we briefly discuss the identification of high $p_T$ charged pion and the techniques used to extract the jet properties.

2. Analysis

2.1. $\pi^\pm$ identification

PHENIX identifies high momentum charged pions with the RICH and EMCal detectors. Charged particles with velocities above the Cherenkov threshold of $\gamma_{th} = 35$ (CO$_2$ radiator) emit Cherenkov photons, which are detected by photo-multiplier tubes (PMTs) in the RICH [7]. This threshold corresponds to 18 MeV/c for electrons, 3.5 GeV/c for muons and 4.9 GeV/c for charged pions. In a previous PHENIX publication [8], we have shown that charged particles with reconstructed $p_T$ above 4.9 GeV/c, which have an associated hit in the RICH, are dominantly charged pions and background electrons from photon conversions. The efficiency for detecting charged pions rises quickly past 4.9 GeV/c, reaching an efficiency of $> 90\%$ at $p_T > 6$ GeV/c.

To reject the conversion backgrounds in the pion candidates, the shower information at the EMC is used. Since most of the background electrons are genuine low $p_T$, they can be rejected by requiring a large shower energy in the EMC. In this analysis, a momentum dependent energy cut at EMC is applied: $E > 0.3 + 0.15p_T$. Additional electron rejection comes from the $\chi^2$ variable,

$$\chi^2 = \sum_i \frac{(E_i^{\text{meas}} - E_i^{\text{pred}})^2}{\sigma_i^2}$$

(2)

where $E_i^{\text{meas}}$ is the energy measured at tower $i$, $E_i^{\text{pred}}$ is the predicted energy for an electromagnetic particle of total energy $\sum_i E_i^{\text{meas}}$ and $\sigma_i$ is the standard deviation for $E_i^{\text{pred}}$. 
Both $E_i^{\text{pred}}$ and $\sigma_i$ are obtained from the electron test beam data. EM shower is more compact than hadronic shower, thus has a smaller $\chi^2$ value. The $\chi^2$ value is then mapped to the probability ($\text{prob}$) for a shower being an EM shower. $\text{prob}$ ranges from 0 to 1, with a flat distribution expected for EM showers and a distribution peaked around 0 for hadronic showers. Figure 2a shows the normalized $\text{prob}$ distribution for the pion candidates and electrons. A cut of $\text{prob} < 0.2$ selects pions with an efficiency of $\gtrsim 80\%$. Since we are interested in per-triggered yield, the detailed knowledge of the pion efficiency is not necessary. The raw pion spectra for requiring only RICH cut and both RICH and EMCal cuts are shown in Figure 2b, the difference between the two is mostly due to electron background. The sample of charged pion used in the correlation analysis is from 5 to 16 GeV/c, with an purity better than 95%.

![Figure 2.](image)

2.2. Extracting jet properties

The correlation functions are generally defined as

$$C(\Delta \phi) = \frac{N_{\text{cor}}(\Delta \phi)}{N_{\text{mix}}(\Delta \phi)} \text{ in azimuth and } C(\Delta \eta) = \frac{N_{\text{cor}}(\Delta \eta)}{N_{\text{mix}}(\Delta \eta)}$$

in pseudorapidity. $N_{\text{cor}}$ and $N_{\text{mix}}$ represent the same-event pair distribution and mixed-event pair distributions, respectively. The mixing is done by pairing trigger $\pi^\pm$ with charged hadrons from events having similar collision vertex and centrality as the $\pi^\pm$. Ref. 2 has shown that the correlation function and the conditional yield are related to each other by just a constant,

$$\frac{1}{N_{\text{trig}}^0} \frac{dN_0^0}{d\Delta \phi} = \frac{R_{\Delta \eta}}{N_{\text{trig}}\epsilon} \frac{N_{\text{cor}}(\Delta \phi)}{2\pi N_{\text{mix}}(\Delta \phi)}$$

where $N_{\text{trig}}^0$ and $N_{\text{trig}}$ are the true and detected number of triggers respectively, and $\epsilon$ is the average single particle efficiency for the associated particles in $2\pi$ in azimuth and $\pm 0.35$ in pseudo-rapidity. $R_{\Delta \eta}$ accounts for the loss of jet pairs outside PHENIX pair acceptance of $|\Delta \eta| < 0.7$.

The two gauss functions in Eq.1 describe the $\Delta \phi$ distribution of the jet signal. The jet signal can also be presented in any other pair variables, such as $\Delta \eta$, trigger $p_T$ ($p_T,\text{trig}$), associated hadron $p_T$ ($p_T,\text{assoc}$), $p_{\text{out}} = p_T,\text{assoc} \sin \Delta \phi$, $x_E = \frac{p_T,\text{assoc}, \cos \Delta \phi}{p_T,\text{trig}}$, di-hadron mass and di-hadron
For every pair variable, we use a statistical weighting method to account for the acceptance correction. According to Eq. 4, each pair on average needs a $\Delta \phi$ dependent correction factor, $w(\Delta \phi)$,

$$w(\Delta \phi) = \frac{R_{\Delta \eta}}{N_{\text{trig}}^{\epsilon}} \frac{1}{2\pi N_{\text{mix}}(\Delta \phi)} \int d\Delta \phi N_{\text{mix}}(\Delta \phi)$$

(5)

When this factor is used as the weight in filling the $x_E$ histograms for both real and mixed pairs, we obtain

$$\frac{1}{N_{\text{trig}}^{\epsilon}} \frac{dN_0}{dx_E} = \sum_{\text{real}} \delta(x_E) w(\Delta \phi)$$

(6)

for the same-event pair distribution. Thus according to Eq. 1, $x_E$ distribution for jet pairs equals to

$$\frac{1}{N_{\text{trig}}^{\epsilon}} \frac{dN^{\text{jet}0}}{dx_E} = \sum_{\text{real}} \delta(x_E) w(\Delta \phi) - C \sum_{\text{mix}} \delta(x_E) w(\Delta \phi)$$

(7)

where

$$C = \frac{BR_{\Delta \eta}}{N_{\text{trig}}^{\epsilon}} \frac{2\pi}{\int d\Delta \phi N_{\text{mix}}(\Delta \phi)}$$

(8)

Replacing $x_E$ with any pair variables, we obtain other jet pair distributions. However, the integral of the jet yield should be conserved independent of the pair variable used, i.e.:

$$\int d\Delta \phi \frac{dN^{\text{jet}0}}{dx_E} = \int dx_E \frac{dN^{\text{jet}0}}{dx_E} = \int dp_{T,\text{assoc}} \frac{dN^{\text{jet}0}}{dp_{T,\text{assoc}}} = \int dp_{out} \frac{dN^{\text{jet}0}}{dp_{out}}$$

(9)

3. Results

3.1. Jet shape

In the following discussion, the trigger $\pi^{\pm} p_T$ is always from 5 to 10 GeV/c, unless specified otherwise. Figure 3 shows the $\pi^{\pm} - h^{\pm} \Delta \phi$ distribution from $p + p$ and $d + Au$ collisions for several range of $p_{T,\text{assoc}}$. The widths decrease with increasing $p_{T,\text{assoc}}$, which is consistent with narrowing of the jet cone for larger $p_{T,\text{assoc}}$. It is interesting to notice that a large fraction of all hadrons in the event are associated with the trigger, thus are originated from the hard-scattered partons. Even for $p_{T,\text{assoc}}$ as low as 0.4 − 1 GeV/c, about 51% hadron yield in $p + p$ (27% in $d + Au$) comes from the jet fragmentation.

Using the event mixing technique, we also measure the jet shape in $\eta$. For the single acceptance of $|\eta| < 0.35$, the pair acceptance in pseudorapidity is limited to be $|\Delta \eta| < 0.7$. Figure 4 shows the same event and mixed event $\Delta \eta$ distribution for $1 < p_{T,\text{assoc}} < 2$ GeV/c, where a cut of $|\Delta \phi| < 1$ is used to select only same side jet pairs. The mixed-event pair distribution is not a perfect triangle due to a gap around $\eta = 0$ in PHENIX central arm detectors. The ratio of the two distributions gives the jet shape in $\Delta \eta$. It is shown in Figure 4 and compared with the jet shape in $\Delta \phi$. There is no significant difference between the two and the extracted widths are consistent in both directions. We extend this comparison to other associated hadron $p_T$ ranges and summarize the results in Figure 5. The overall agreement between the jet widths in $\Delta \eta$ and $\Delta \phi$ is pretty good, except at small $p_{T,\text{assoc}}$, where the width in $\Delta \eta$ is systematically lower than that in $\Delta \phi$. The fact that this discrepancy exist in both $p + p$
and d + Au collisions indicates that this deviation is likely due to the systematics of the fitting in a limited Δη range rather than any real physics effect in d + Au.

From the measured jet width, one can calculate the rms value of $z k_{T_y}$ [2], $(z k_{T_y})_{RMS} = \sqrt{\langle z^2 k_{T_y}^2 \rangle}$, for both $p+p$ and d + Au. The resulting $(z k_{T_y})_{RMS}$ is plotted as function of trigger $p_T$ in Figure 6. It looks quite similar between $p+p$ and d + Au. The difference of $(z k_{T_y})_{RMS}$, averaged over $p_T$, is $\langle z^2 k_{T_y}^2 \rangle_{dAu} - \langle z^2 k_{T_y}^2 \rangle_{pp} = 0.64 \pm 0.78 \pm 0.42$ (GeV/c)$^2$, which is consistent with 0. According to various theoretical estimations [3], the typical contribution to $\langle k_{T_y}^2 \rangle$ from multiple scattering is 1 (GeV/c)$^2$ in central d + Au collisions, while the contribution from initial and final radiation is much larger (around 8 (GeV/c)$^2$) [10]. The small multiple scattering

Figure 3. Corrected conditional pair distributions for $p+p$ and minimum bias d+Au collisions. The trigger $\pi^{\pm}$ are correlated with hadrons with $p_{T,assoc}$ 0.4–1.0 GeV/c, 1.0–2.0 GeV/c, 2.0–3.0 GeV/c and 3.0–5.0 GeV/c (from top to bottom and left to right).

Figure 4. a) The same-event and mixed-event pair distribution in Δη, b) the correlation function in Δη (open boxes) and Δφ (filled circles).
contribution might have been washed out by the much larger contributions from initial and final radiation, which explains the lack of difference between the two systems.

An alternative but more direct way in studying the away side $k_T$ broadening is through the $p_{out}$ distribution. For small angles, $p_{out}$ has simple relation to $j_T$ and $k_T$ [2]:

$$\langle p_{out,\text{same}}^2 \rangle \approx \langle j_T^2 \rangle + x_E^2 \langle j_T^2 \rangle$$

$$\langle p_{out,\text{away}}^2 \rangle \approx \langle j_T^2 \rangle + x_E^2 \langle j_T^2 \rangle + 2 x_E^2 \langle z^2 k_T^2 \rangle$$

Thus the difference of same side and away side $p_{out}$ distributions directly reflects the contribution from $k_T$.

$$2 x_E^2 \langle z^2 k_T^2 \rangle \approx \langle p_{out,\text{same}}^2 \rangle - \langle p_{out,\text{away}}^2 \rangle$$

Figure 7a shows the same side and away side $p_{out}$ distribution. There is a significant difference between the two, which reflects the contribution from $k_T$. Both distributions have a gauss shape at small $p_{out}$ followed by an excess at large $p_{out}$. The gauss part presumably is due to the jet fragmentation (in both the same and away side) and intrinsic $k_T$ (away side only), while the excess is evidence for hard radiation contribution of the outgoing partons. Since the away side
$p_{\text{out}}$ carries information about $k_T$, we compare between $p+p$ and $d+Au$ to see whether there is hint of additional $k_T$ broadening in $d+Au$. The comparisons of away side $p_{\text{out}}$ distributions are shown in Figure 7, no apparent differences are observed, consistent with the observations that $(z k_{T_y})_{\text{RMS}}$ are similar between $p+p$ and $d+Au$.

![Figure 7.](attachment:image7.png)

**Figure 7.** a) Same side and away side $p_{\text{out}}$ distributions in $d+Au$ collisions. b) The away side $p_{\text{out}}$ distributions compared between $p+p$ and $d+Au$.

### 3.2. Jet yields

The same side and away side $p_T$ distributions of the charged hadrons associated with jets are plotted in Figure 8, comparing between $p+p$ and $d+Au$ collisions. The same side yield is related to the di-hadron fragmentation, since both particles comes from the same jet, while the away side yield depends on two independent fragmentation functions: one parton fragments to produce the trigger, while the second parton produces the associated hadron. No apparent differences are seen between $p+p$ and $d+Au$; this observation is in contradiction to some recombination model prediction [11], in which a significant difference is expected due to shower-thermal contribution.

![Figure 8.](attachment:image8.png)

**Figure 8.** Jet pair distribution as function of $p_{T,\text{assoc}}$ for same side (right panel) and away side (left panel) in $p+p$ and $d+Au$.  

Di-hadron correlation also gives the $x_E$ distribution $\frac{1}{N_{trig}} \frac{dN_h}{dx_E}$, where $x_E = z_{assoc}/z_{trig}$. When di-jet $p_T$ imbalance is ignored and $z_{assoc} \ll z_{trig}$, $z_{trig}$ varies very slowly with $z_{assoc}$. Hence the $x_E$ distribution is closely related to the fragmentation function $D(z)$,

$$
\frac{1}{N_{trig}} \frac{dN_h}{dx_E} \approx z_{trig} D(z)
$$

(11)

Figure 9 shows the measured $x_E$ distribution between $p+p$ and $d+Au$. Again, no difference is seen between the two collision systems in both the same side and away side.

![Figure 9](image)

**Figure 9.** $x_E$ distributions as function of $p_T,assoc$ for same side (right panel) and away side (left panel) in $p+p$ and $d+Au$.

In $e^+e^-$ or $p+p$ collisions, the fragmentation functions $D(z)$ are known to approximately scale, i.e. are independent of jet energy. To check whether this is still true in $d+Au$ collisions, we plot in Figure 10 the conditional yields as function of trigger $p_T$ in different ranges of $x_E$ for both $p+p$ and $d+Au$. The amount of variation is within $\pm 25\%$ for $p_T$ from $5-12$ GeV$/c$ with very little difference between the two systems. So we conclude that the evolution of the jet fragmentation as function of jet energy is very similar between $p+p$ and $d+Au$.

### 3.3 Comments on underlying event

Events triggered by high $p_T$ hadrons not only contain particles originated from the two hard-scattered partons, but also those come from soft multiple interaction and the beam remnants. Underlying event in $p+p$ and $d+Au$ collisions refers to all hadrons except those from the two outgoing hard-scattered partons, which includes contributions from the beam remnants and initial and final state radiation [12]. The physics of the underlying event is poorly known due to its non-perturbative nature. It is often studied phenomenologically with various QCD Monte-carlo models that have been tuned to fit the data [13]. Underlying event has been studied extensively at the Tevatron energy [12, 14]. Similar studies at the RHIC are very useful in understanding it’s dependence on $\sqrt{s}$, and can provide valuable constrains on the underlying event physics at the LHC.

Figure 11 shows the jet pair distribution in $p+p$ collisions, reproduced from Figure 3. The pedestal in the $\Delta \phi$ correlation, which represents the underlying event contribution, decreases quickly and becomes negligible at $p_{T,assoc} > 2$ GeV$/c$. However, the level corresponding to minimum bias $p+p$ events, denoted by the thick horizontal line, seems to decrease even faster.
Figure 10. Away side conditional yield as function of $p_{T,\text{trig}}$ for different ranges of $x_E$ in a) minimum-bias $d + Au$ collisions and b) $p + p$ collisions.

Since minimum bias event has small hard-scattering contribution, the relative abundance of the pedestal in triggered events over the minimum bias events indicates that most of the underlying event comes from the initial or final state radiation of the hard-scattered partons.

Figure 11. Corrected condition yield in $\Delta\phi$ for $p + p$ collisions (from Figure 3). The thick solid line represents the average level for minimum bias events, i.e. it is equal to $\text{Yield}_{pp}/(2\pi)$.

The underlying event at RHIC can be checked in QCD Monte-carlo models. We use the PYTHIA6.131, which seems to be able to reproduce the jet conditional yield as shown in Figure 12. The roles of the initial/final state radiations are studied by switching them on and
off in PYTHIA simulation. Figure 13 shows a typical $\pi^\pm - h^\pm$ azimuthal correlation with (top histogram) and without (bottom histogram) radiation from the simulation. There is a significant enhancement in the pedestal level when radiations are enabled. We can perform a more quantitative comparison by plotting separately the jet contribution (double gauss component), the underlying event (the constant component) and the minimum bias event level as function of $p_T$ in Figure 13. The hierarchy of the three contributions can be clearly seen. For event tagged with a high $p_T$ jet, the spectra for both the jet fragmentation and the underlying events are much harder than that from typical minimum bias events. Current statistics from $p + p$ does not allow a quantitative comparison with the models yet, a much larger dataset collected from recent RHIC $p + p$ run in 2005 should help to address this question in the near future.

![Figure 12](image1)

**Figure 12.** Conditional yield with trigger pion in $5 < p_{T,\text{trig}} < 6$ GeV/c, compared with PYTHIA 6.131 for both same side (left panel) and away side (right panel).

![Figure 13](image2)

**Figure 13.** PYTHIA MC simulation: a) $\Delta \phi$ distribution for $3 < p_{T,\text{assoc}} < 5$ with (top histogram) and without (bottom histogram) radiation. b) the $p_T$ spectra for jet pairs and underlying event in $p + p$ events with a $> 5$ GeV/$c$ trigger, compared with minimum bias event yield.

What about the underlying event in $d + Au$? Figure 13 indicates that the underlying event levels are larger than those in $p + p$, although the properties of the jets are quite similar. Since the hard-scattering only happens in one of the nucleon-nucleon collision in $d + Au$, we can
assume the ambient particle production mechanism is the same as in minimum bias nucleon-nucleon collision. In this case, the ambient particle production should simply scale as the nuclear modification factor, \( R_{dAu} \) measured in \( d + Au \) \[15\]. Thus the underlying event yields in \( p + p \) and \( d + Au \), \( U_{dAu} \) and \( U_{pp} \) are connected to each other through the following simple relation,

\[
U_{dAu} = U_{pp} + R_{dAu} (N_{coll} - 1) \text{Yield}_{pp}
\] (12)

where \( \text{Yield}_{pp} \) represents the hadron yield per event in minimum bias \( p + p \) collisions. Divide both side by \( \text{Yield}_{pp} \), we get,

\[
\lambda_{dAu} = \lambda_{pp} + R_{dAu} (N_{coll} - 1)
\] (13)

\[
\lambda_{dAu} = U_{dAu}/\text{Yield}_{pp}, \lambda_{pp} = U_{pp}/\text{Yield}_{pp}
\] (14)

note \( \lambda_{pp} \) denotes the ratio of underlying event yield to minimum bias event in \( p + p \), which should be larger than 1 according to Figure.11 and Figure.13.

4. Conclusion
The di-jet decay kinematics are studied using \( \pi^\pm - h^\pm \) correlation in \( p + p \) and \( d + Au \) collisions. Measured jet widths, the calculated \( k_T \) and distributions of \( p_{out} \) are very similar between \( p + p \) and \( d + Au \), which indicate no or small broadening in cold nuclear medium. Jet yield distribution in associated hadron \( p_T \) and \( x_E \) are also similar between \( p + p \) and \( d + Au \), consistent with no significant increase in jet multiplicity in \( d + Au \) relative to \( p + p \). The dependence of the \( x_E \) distribution on trigger \( p_T \) is weak in measured trigger \( p_T \) range. The underlying event yield in \( p + p \) is studied in PYTHIA Monte-carlo. Events containing a large \( p_T \) trigger appear to have an underlying event spectra much harder than the minimum bias hadron spectra.

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