Charged Higgs in the NMSSM

A. G. Akeroyd
Department of Physics, National Central University, Cheungli, Taiwan
E-mail: akeroyd@ncu.edu.tw

Abdesslam Arhrib
Département de Mathématiques, Faculté des Sciences et Techniques, B.P 416 Tangier, Morocco,
LPHE, Département de Physique, Faculté des Sciences, Marrakesh, Morocco
E-mail: aarhrib@ictp.it

Qi-Shu Yan
Physics Department, University of Toronto, 60 St George Street, Toronto,
Ontario Canada M5S 117
E-mail: yangs.yan@utoronto.ca

The charged Higgs boson decays $H \rightarrow W A_1$ and $H \rightarrow W h_1$ are studied in the framework of the next-to Minimal Supersymmetric Standard Model (NMSSM). It is found that the decay rate for $H \rightarrow W A_1$ can dominate both below and above the top-bottom threshold. We suggest that $pp \rightarrow H A_1$ is a promising discovery channel for a light charged Higgs boson in the NMSSM with small or moderate $\tan \beta$ and dominant decay mode $H \rightarrow W A_1$. This $W A_1A_1$ signature can also arise from the Higgsstrahlung process $pp \rightarrow h_1$ followed by the decay $h_1 \rightarrow A_1A_1$. It is shown that there exist regions of parameter space where these processes can have comparable cross sections and we suggest that their respective signals can be distinguished at the LHC by using appropriate reconstruction methods.

Prospects for Charged Higgs Discovery at Colliders
September 16-19 2008
Uppsala, Sweden
1. A charged Higgs boson (H) appears in any extension of the Standard Model with two hypercharge Y=1 doublets. Its phenomenology has been extensively studied in both the Two Higgs Doublet Model (2HDM) and MSSM. The presence of H is also predicted in the Next-to MSSM (NMSSM) in which an additional singlet neutral complex scalar field S is added to the two Higgs doublets of the MSSM.

In the NMSSM, after electroweak symmetry breaking the Higgs spectrum consists of three neutral scalars (h_1, h_2, h_3), two pseudoscalars (A_1, A_2) and a pair of charged Higgs bosons H. In both the CP-odd and CP-even sector the physical eigenstates are ordered as M_{h_1} < M_{h_2} < M_{h_3} and M_{A_1} < M_{A_2}. For detailed discussions of the Higgs sector of the NMSSM the reader is referred to [1, 2, 3, 4, 5]. The mass of H at tree-level is given by [6, 7]:

$$M_H^2 = \frac{2\mu_{eff}}{\sin 2\beta} (\lambda_S + \kappa S) + M_W^2 - \lambda^2 v^2$$  \hspace{1cm} (1)$$

where tan\beta = v_u/v_d and v^2 = v_u^2 + v_d^2. This differs from the corresponding MSSM expression in which M_A and M_H are strongly correlated and become roughly equal for M_A \approx 140 \text{ GeV}.

The CP-odd mass matrix can be obtained as follows: Firstly, as in MSSM one rotates the bare fields (\tilde{\gamma}_m H_u, \tilde{\gamma}_m H_d, \tilde{\gamma}_m S) into a basis (\gamma G, \beta m S) where G is a massless Goldstone boson. Then one eliminates the Goldstone mode and the remaining 2 CP-odd states are:

$$A_1 = \cos \theta_A \alpha + \sin \theta_A \beta \gamma (S) \hspace{1cm} ; \hspace{1cm} A_2 = \sin \theta A + \cos \theta_A \beta \gamma (S)$$  \hspace{1cm} (2)$$

Where A = cos\beta \gamma \gamma_m H_u + sin\beta \gamma \gamma_m H_d is the CP-odd MSSM Higgs boson while \beta \gamma m comes from the singlet S field.

In the MSSM the coupling \text{H} \text{AW} (where A is the CP-odd neutral Higgs boson) contains no mixing angle suppression but the relation M_A \approx M_H ensures that the decay \text{H} \text{AW} is greatly suppressed in most of the parameter space. In the NMSSM, the relevant couplings for our study are described by the following Lagrangian:

$$\mathcal{L}_{VVH/HH} = g_{WW} g_{VVH} W^+ \mu W_\mu h_i + g_{WW} \left( \frac{ig_{W^+ H}}{2} h_i + \frac{P_{ij}}{2} A_1 \right) \frac{\partial^\mu H}{\partial \nu}$$  \hspace{1cm} (3)$$

where g_{VVH} = \sin \beta S_1 + \cos \beta S_2, g_{W^+ H} h_i = \cos \beta S_2, P_{ij} = \cos \theta_A and P_{ij} = \sin \theta_A, S and P are orthogonal matrices which diagonalize respectively the CP-even and CP-odd scalar mass matrix. From the last term in eq. (3) one can see that the vertex W \text{H} A_1 is directly proportional to P_{11} i.e. the doublet component of the mass eigenstate A_1. Consequently, if A_1 is entirely composed of doublet fields this coupling is maximized and if A_1 is purely singlet the coupling vanishes.

Now we are ready to describe the phenomenology of the H in the NMSSM and we summarize the results of our earlier work [7]. The phenomenology of H in the NMSSM has many similarities with that of H in the MSSM. This is to be expected since the fermionic couplings are identical in the two models. The main differences in their phenomenology originate from the possibility of large mass splittings among the Higgs bosons in the NMSSM which permits decay channels like \text{H} \text{AW} to proceed on-shell [8]. Moreover, in the NMSSM a light CP-even h_1 is also allowed and one can have the opening of the decay \text{H} \text{h_1 W} both below and above the top-bottom threshold. This latter channel may change the NMSSM phenomenological predictions for the charged Higgs with respect to the MSSM [8]. In the MSSM the decay \text{H} \text{h_1 W} is also open but the coupling g_{W^+ H} h_i \cos^2 (\beta \alpha) is strongly suppressed when M_H \approx m_{h_1} + m_w and thus its branching ratio is very small for such M_H. For M_H < m_{h_1} + m_w and just above the threshold the branching ratio for this channel can reach 10% at most for small values of tan\beta [9, 10, 11].
The decay $H \rightarrow AW$, where $A$ is a CP-odd Higgs boson, may be sizeable in a variety of models with a non-minimal Higgs sector such as Two Higgs doublet models (Type I and II) \cite{12,13,14} and in SUSY models with Higgs triplets \cite{15}. Two LEP collaborations (OPAL and DELPHI) performed a search for a charged Higgs decaying to $AW$ (assuming $m_A > 2m_b$) and derived limits on the charged Higgs mass \cite{16} comparable to those obtained from the search for $H \rightarrow cs \tau v$. In the MSSM the decay width for $H \rightarrow AW$ is very suppressed in most of the parameter space \cite{8,10} because the charged Higgs and the CP-odd Higgs are close to mass degeneracy. The importance of the decays $H \rightarrow A_1 W$ and $H \rightarrow h_1 W$ in the NMSSM was first pointed out in \cite{8}. Their branching ratios may be close to 100% which can provide a clear signal at the LHC.

The decay width of $H \rightarrow A_1 W$ is directly proportional to $\cos \theta_A$ which is the doublet component of $A_1$. This decay width can be substantially enhanced if $A_1$ is predominantly composed of doublet fields. However, even with small doublet (large singlet) component of $A_1$ it is possible that $H \rightarrow A_1 W$ is the dominant decay mode. We perform a scan of the parameter space using the code NMSSM-Tools \cite{17} in order to quantify the importance of $H \rightarrow A_1 W$ and $H \rightarrow h_1 W$.

Hereafter we assume that all scalar superparticles share the same soft mass term $M_{SUSY}$ and the ratios of gaugino masses satisfy $M_1 : M_2 : M_3 = 1 : 2 : 6$; the trilinear couplings are related to $M_{SUSY}$ but the sign is not fixed, i.e. $A_1, A_2 = 2M_{SUSY}$. We scan the parameter space of the model by varying the free parameters within the following region:

$$\lambda = [0; 1]; \kappa = [1; 1]; \tan \beta = [0.2; 60]; \mu = [1; 1] \text{TeV};$$

$$A_1 = [1; 0; 0] \text{TeV}; A_2 = [1; 0; 0] \text{TeV}; M_{SUSY} = [0; 2; 3] \text{TeV}; M_1 = [0; 0; 0; 3] \text{TeV};$$

While varying these parameters, we take into account the experimental constraints on the MSSM spectrum e.g., charged Higgs mass $\geq 80$ GeV, chargino and scalar fermions $\geq 100$ GeV. We also apply the full set of LEP constraints obtained from searches for neutral Higgs bosons decaying to final states like $Z2b, Z4b, 6b, 6\tau, Z2b2\tau, Z4\tau, 2b2\tau$.

In Fig. (1) we display the branching ratios of $W A_1, \tau v$ and top-bottom modes. Before the opening of the $H \rightarrow tb$ channel, the full dominance of $W A_1$ over $\tau v$ requires light $M_{A_1} < 100 GeV$, large doublet component of $A_1$ and $\tan \beta$ not too large. Note that at large $\tan \beta \approx 15$ $25$, the $W A_1$ and $\tau v$ channels become comparable in size. Once the decay $H \rightarrow tb$ is open, it competes strongly with $W A_1$ for $\tan \beta < 15$. As can be seen from Fig. (1) left, the branching
ratio of $H \, W \, A_1$ is less than 90%. It is interesting to see also that for $\cos^2 \theta_A < 0.05$ there is not a single point with $\text{Br}(H \, W \, A_1) > 0.5\%$. Note also that at large $\tan \beta > 25$, it is hard for $H \, W \, A_1$ to compete with $\tau \nu$ and top-bottom modes.

The most problematic region for $H$ discovery in the MSSM is for moderate values of $\tan \beta$, since the production mechanisms which rely on a large bottom quark or top quark Yukawa coupling (e.g. $gb \cdot H \, t$) are least effective. Hence alternative mechanisms which could offer good detection prospects for $H$ at moderate values of $\tan \beta$ are desirable. The cross sections for the pair production mechanisms $pp \rightarrow H \, A_1$ and $pp \rightarrow H \, h_1$ fall quickly with increasing scalar masses but for relatively light masses ($< 200$ GeV) they can provide promising signal rates which might enable their detection at the LHC (see [18] for studies in the context of the MSSM). One common feature is that the produced scalars enjoy large transverse momenta, which are crucial for the trigger and event selection.

In the NMSSM, if the coupling $H \, W \, A_1$ is sizeable, so will be the cross section for $pp \rightarrow H \, A_1$ provided that $H$ and $A_1$ are not too heavy. The production mechanism $pp \rightarrow H \, A_1$ followed by the decay $H \rightarrow W \, A_1$ would give rise to a signal $W \, A_1 A_1 \rightarrow Wbb \nu \nu$ or $W \, A_1 A_1 \rightarrow W \tau \tau \tau \tau$. The signature $W \, A_1 A_1 \rightarrow Wbb \nu \nu$ was simulated at the LHC in [20] in the context of the CP violating MSSM with the conclusion that a sizeable signal essentially free of background could be obtained. We use NMSSM-TOOLS1.1.1 to calculate the mass spectrum and couplings of the NMSSM Higgs bosons, and we link CTQ6.1M PDF distribution to this code in order to calculate the cross sections of $pp \rightarrow H \, A_1$, $pp \rightarrow H \, h_1$ and $pp \rightarrow W \, h_1$. All cross sections are evaluated at a scale which is the sum of the masses in the final states and do not include next-to-leading order QCD enhancement factors (K factors) of around 1.2–1.3 [18], [21].

Note that the process $pp \rightarrow H \, A_1$ leads to the same signature as the process $pp \rightarrow V h_1$ and $W A_1 A_1 \rightarrow Wbb \nu \nu$. The latter has been simulated in [22] and also offers very good detection prospects. We will compare the magnitude of these two distinct mechanisms which lead to the same $Wbb \nu \nu$ signature. In addition, the mechanism $pp \rightarrow H \, h_1$ followed by the decay $H \rightarrow W \, A_1$, would also lead to the same final state $W \, A_1 h_1 \rightarrow Wbb \nu \nu$.

Hence a numerical comparison of their cross sections is of particular interest and is shown in Fig. (3), where all points satisfy the following conditions:

$$\sigma(pp \rightarrow H \, A_1) > 0.1 \text{ pb} \quad \text{and} \quad \sigma(pp \rightarrow W \, h_1) > 0.1 \text{ pb} :$$ (5)

Superimposed on Fig. (2a) and Fig. (2b) are the main decay modes of the charged Higgs boson and the decay neutral Higgs boson $H_1$ respectively. We further impose the following conditions:

$$\text{Br}(H \rightarrow W \, A_1) > 0.5 \quad \text{and} \quad \text{Br}(h_1 \rightarrow A_1 A_1) > 0.5 ;$$ (6)

and the surviving points are displayed in Fig. (7a). Importantly, there are many points where the two cross sections are of comparable size. We note that for these points in Fig. (7a) the pseudoscalar $A_1$ can be both $R$-axion like or a mixture of the three allowed basic axions. If the magnitude of the cross sections of both $pp \rightarrow H \, A_1$ and $pp \rightarrow V h_1$ are similar then the interference of the two channels (i.e., the same $Wbb \nu \nu$ signature arising from distinct production mechanisms) should be taken into account. We have neglected such effects in the present study.

We now discuss whether the $Wbb \nu \nu$ signatures can be distinguished experimentally by comparing the strategies adopted in [24] (for $pp \rightarrow H \, A^0$) and [22] (for $pp \rightarrow W \, h_1$). In order to reconstruct the peak of the CP-even Higgs $h_1$, one can select events with a charged lepton and four tagged $b$ quark jets as shown in [22]. This enables both a clean Higgs signal with high significance and a measurement of $M_{h_1}$ given by the invariant mass of the four $b$ quark jets, $m_{4b}$. The process
Charged Higgs in the NMSSM

Abdesslam Arhrib

Figure 2: Left panel: comparison of $\sigma(pp \rightarrow H \ A_1)$ and $\sigma(pp \rightarrow W \ h_1)$ with two $H$ decay modes. Right panel: comparison of $\sigma(pp \rightarrow H \ A_1)$ and $\sigma(pp \rightarrow W \ h_1)$ with two $h_1$ decay modes. The dotted line corresponds to $\sigma(pp \rightarrow W \ h_1) = \sigma(pp \rightarrow H \ A_1)$.

$pp \rightarrow H \ A_1$ might be an irreducible background but presumably could be significantly suppressed with the aforementioned cut on $m_{4b}$ e.g., $m_{h_1} - 15 \text{GeV} < m_{4b} < m_{h_1} + 15 \text{GeV}$.

Regarding detection of $pp \rightarrow H \ A_0^0$, it was demonstrated in [20] (for the analogous process $pp \rightarrow H \ H_1$ in the CP violating MSSM) that the mass of $H$ can be reconstructed. This is achieved by defining a transverse mass ($M_T$) which is a function of the momenta of the two secondary $b$ jets (i.e., those originating from the decay $H \rightarrow A_1 W \rightarrow W bb$) and the momenta of the lepton and missing energy coming from the $W$ boson. It was shown that $M_T$ is sensitive to the underlying charged Higgs mass and thus can be used for the determination of $M_H$. The pair of $b$ jets from $pp \rightarrow W \ h_1$ might be an irreducible background but presumably could be suppressed with a cut on $M_T$.

To reconstruct the peak of the light CP-odd neutral Higgs $A_1$, one can require events with four tagged $b$ jets, construct the three possible double pairings of $b\bar{b}$ invariant masses, and then select the pairing giving the least difference between the two $b\bar{b}$ invariant masses values [21]. $W4b$ signatures from the process $pp \rightarrow W \ h_1$ also contribute constructively to the reconstruction of $A_1$. Thus we conclude that it is promising to reconstruct the peaks of the CP-even neutral Higgs ($h_1$), charged Higgs ($H^\pm$) and CP-odd neutral Higgs ($A_1$) and thus experimentally distinguish the $Wbbb$ signatures arising from the two distinct production mechanisms. We defer a detailed simulation to a future work.

In summary, it was shown that $H \rightarrow W A_1$ can dominate over the standard decays $H \rightarrow \tau \nu$ and $H \rightarrow tb$ both below and above the top-bottom threshold. Large branching ratios for $H \rightarrow W A_1$ and $H \rightarrow W h_1$ would affect the anticipated search potential for $H$ at the LHC. We also studied the production process $pp \rightarrow H A_1$ and showed that sizeable cross sections ($>1$ pb) are possible. It is known that intermediate values of $\tan\beta$ (e.g., $5 < \tan\beta < 20$) are most problematic for discovery of $H$ at the LHC [23], since the $H \rightarrow tb$ Yukawa coupling (which is
employed in the conventional production processes) takes its lowest values. In such a region the process $pp \rightarrow H A_1$ can have a sizeable cross section if $m_H + m_{A_1} < 200$ GeV. Therefore we propose $pp \rightarrow H A_1$ as a unique mechanism to probe the parameter space of intermediate $\tan\beta$ and light charged Higgs boson in the NMSSM.

**Acknowledgments**

A.A thanks the organiser for the warm hospitality extended to him during the workshop.

**References**

[1] M. Drees, Int. J. Mod. Phys. A **4**, 3635 (1989).
[2] T. Elliott, S. F. King and P. L. White, Phys. Rev. D **49**, 2435 (1994).
[3] F. Franke and H. Fraas, Int. J. Mod. Phys. A **12**, 479 (1997).
[4] D. J. Miller, R. Nezvorov and P. M. Zerwas, Nucl. Phys. B **681**, 3 (2004).
[5] Chapter 4 in E. Accomando *et al.*, “Workshop on CP studies and non-standard Higgs physics,” arXiv:hep-ph/0608079.
[6] S. F. King and P. L. White, Phys. Rev. D **53**, 4049 (1996).
[7] A. G. Akeroyd, A. Arhrib and Q. S. Yan, Eur. Phys. J. C **55**, 653 (2008).
[8] M. Drees, E. Ma, P. N. Pandita, D. P. Roy and S. K. Vempati, Phys. Lett. B **433**, 346 (1998).
[9] S. Moretti and W. J. Stirling, Phys. Lett. B **347**, 291 (1995) [Erratum-ibid. B **366**, 451 (1996)].
[10] A. Djouadi, J. Kalinowski and P. M. Zerwas, Z. Phys. C **70**, 435 (1996).
[11] M. Drees, M. Guchait and D. P. Roy, Phys. Lett. B **471**, 39 (1999).
[12] F. Borzumati and A. Djouadi, Phys. Lett. B **549** (2002) 170.
[13] A. G. Akeroyd, Nucl. Phys. B **544**, 557 (1999).
[14] A. G. Akeroyd, A. Arhrib and E. Naimi, Eur. Phys. J. C **20**, 51 (2001); Eur. Phys. J. C **12** (2000) 451 [Erratum-ibid. C **14** (2000) 371].
[15] J. L. Diaz-Cruz, J. Hernandez-Sanchez, S. Moretti and A. Rosado, Phys. Rev. D **77**, 035007 (2008).
[16] J. Abdallah *et al.* [DELPHI Collaboration], Eur. Phys. J. C **34** (2004) 399; OPAL Phys. Note PN445 (2000); OPAL Phys. Note PN472 (2001).
[17] U. Ellwanger, J. F. Gunion and C. Hugonie, JHEP **0502**, 066 (2005); NMSSMTOOLS link can be found @ http://www.th.u-psud.fr/NMHDECAY/nmssmt.jpg.html.
[18] S. Kanemura and C. P. Yuan, Phys. Lett. B **530**, 188 (2002); Q. H. Cao, S. Kanemura and C. P. Yuan, Phys. Rev. D **69**, 075008 (2004).
[19] A. G. Akeroyd, Phys. Rev. D **68**, 077701 (2003).
[20] D. K. Ghosh and S. Moretti, Eur. Phys. J. C **42**, 341 (2005).
[21] T. Han and S. Willenbrock, Phys. Lett. B **273**, 167 (1991); M. Spira, Fortsch. Phys. **46**, 203 (1998).
[22] K. Cheung, J. Song and Q. S. Yan, Phys. Rev. Lett. **99**, 031801 (2007).
[23] K. A. Assamagan, Y. Coadou and A. Deandrea, Eur. Phys. J. direct C **4**, 9 (2002).