To obtain the interactions which cause neutrino flavor conversion, we introduce a horizontal symmetry into the standard model (SM) and propose the hypothesis that new interactions generated by the horizontal symmetry lead to neutrino flavor conversion and oscillation. To support our hypothesis, we evaluate the flavor conversion probability by new interactions by utilizing the definition of cross section, and the prediction is consistent to experimental data. From our hypothesis, neutrino oscillation is fluctuation of flavor distribution before arriving at equilibrium.

**Keywords:** neutrino-neutrino interaction, neutrino oscillation, beyond SM.

**Introduction**

The phase evolution can depict phenomenologically neutrino oscillation [1] very well, and also manifest neutrinos massive. But the interactions which cause neutrino flavor conversion remain puzzling [2,3], which we will attempt to investigate in this paper.

No interactions implied in the SM can induce neutrino flavor conversion, while a horizontal symmetry added into SM will provide these interactions [4] and also the masses of neutrinos [5]. Therefore we extend the SM and propose a hypothesis that new interactions from horizontal symmetry lead to neutrino flavor conversions. And neutrino oscillation is the macro phenomenon of the
sum of all interactions before equilibrium, which will results in a mixing matrix of anarchy [6,7].

To support our hypothesis, we will calculate the flavor conversion probability by new interactions and compare it to the experimental data. In these calculation, we will assume neutrinos collide due to ‘Brown movement’ with their moments of all possible directions. We will uncover that it is their mass differences that lead to three different neutrino mixing angles.

This paper will be constructed as follows. We will introduce the model of $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$, which is an extension of SM by adding a horizontal symmetry $SU(2)_N$. We will evaluate the flavor conversion probability induced by new interactions and compare it with experimental data. We will demonstrate the resultant phenomenon of neutrino oscillation.

### The model of $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$

For simplicity, we will take two-flavor frame. Thus, our model contains four leptons $\nu_e, e, \nu_\mu, \mu$ and four quarks with neglecting the color degree of freedom. (For anomaly free, quark sector is assigned the same as SM and we will not discuss them. And we will not discuss the source of neutrino mass in this paper either.)

The total group is $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$. Considering suppression of $e \leftrightarrow \mu$ in experiments, the horizontal flavor symmetry $SU(2)_N$ only works between neutrinos. Thus the up sectors of SM doublets are under the horizontal symmetry $SU(2)_N$, while the lower sectors are not. We will assign the scale which breaking $SU(2)_N$ just a little higher than the scale of SM. After $SU(2)_N$ broken, the model should return to SM.

For $2 \times 2$ representation of field like $\psi$, the transformation under the group $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$ can be equivalent to transformation under group $SU(2)_L, U(1)_Y$ and $SU(2)_N$ in the meantime. (1) is the transformations of $\psi$ under group $SU(2)_L, U(1)_Y$ and $SU(2)_N$ respectively:

$$
\psi \rightarrow e^{-i\frac{1}{2}a(x)\tau_w^a} \psi \\
\psi \rightarrow e^{-i\frac{1}{2}b(x)\hat{Y}} \psi \\
\psi \rightarrow \psi e^{i\frac{1}{2}c(x)\tau_F^b} 
$$

In (1), $\tau_w, \tau_F$ are generators associated with $SU(2)_L$ and $SU(2)_N$ respectively, $\hat{Y}$ is the weak hypercharge operator related to $U(1)_Y$. The covariant derivatives related with transformation of $\psi$ under $SU(2)_L, U(1)_Y$ and $SU(2)_N$ respectively are
\[ D_\mu = \partial_\mu - i \frac{1}{2} g_w \tau^a_i W^a_{\mu} \] (2)

\[ D_\mu = \partial_\mu - i \frac{1}{2} g_Y Y_\mu \hat{Y} \] (3)

\[ D_\mu \psi = \partial_\mu \psi + i \frac{1}{2} g_Y \psi \tau^a_i F^a_{\mu} \] (4)

In our model, two neutrino flavors are assigned into the doublet representation of \( SU(2)_N \) as follows

\[
\begin{pmatrix}
  \nu_e \\
  \nu_\mu
\end{pmatrix}
\quad (2,2,-1)
\] (5)

The numbers within parentheses stand for \( SU(2)_L, SU(2)_N \) and \( U(1)_Y \) quantum numbers \((2I + 1, 2I_N + 1, Y)\) where \( I, I_N \) are isospins of the subgroups \( SU(2)_L, SU(2)_N \) respectively, and \( Y \) is the \( U(1) \)-hypercharge.

The \( SU(2)_L \otimes U(1)_Y \) is the W−S model and the doublets are

\[
\begin{pmatrix}
  \nu_e \\
  e
\end{pmatrix}
_L, \quad \begin{pmatrix}
  \nu_\mu \\
  \mu
\end{pmatrix}_L
\] (6)

and the lower sectors \( e_L, \mu_L \) are assigned as \((2,1,-1)\).

We assign all leptons in a left-handed \( 2 \times 2 \) matrix

\[
\begin{pmatrix}
  \nu_e & \nu_\mu \\
  e & \mu
\end{pmatrix}_L
\] (7)

and two right handed singlets

\[ R_e = e_R, \quad R_\mu = \mu_R \quad (1,1,-2) \] (8)

The electric charge formula is given by

\[ Q = I_{3L} + Y/2 \] (9)

Thus, there are seven vector bosons \( W^i_\alpha, F^i_\alpha \ (i = 1 \sim 3) \) and \( Y_\alpha \), associated with the subgroups \( SU(2)_L, SU(2)_N, U(1)_Y \) respectively.

The fermion dynamical Lagrangian is

\[
\mathcal{L}_f = i \bar{R}_e \gamma^\alpha (\partial_\alpha + i g_Y Y_\alpha) R_e + i \bar{R}_\mu \gamma^\alpha (\partial_\alpha + i g_Y Y_\alpha) R_\mu \\
+ i \text{Tr} [\bar{L} \gamma^\alpha (\partial_\alpha + \frac{1}{2} g_Y Y_\alpha - i \frac{1}{2} g_w \tau^i \omega W^i_{\alpha}) L] \\
+ i \text{Tr} [\bar{L} \gamma^\alpha i \frac{1}{2} g_F \frac{1}{2} (1 + \tau_3) L \tau^i_F (F^i_\alpha)]
\] (10)
The difference of (10) from [4] is that $\tau_i^F$ only act on neutrino fields and vector bosons $F_\alpha^i$ only intermediate neutrino-neutrino interactions, and the term $\frac{1}{2}(1 + \tau_3)$ selects the up sectors of $L$ for horizontal symmetry $SU(2)_N$. The proof of the fermion dynamical Lagrangian (10) invariant under the group $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$ is shown in Appendix A.

Though our horizontal symmetry exists just between neutrinos, the masses of vector bosons are given by Higgs scalars vacuum expectation value (VEV) and the mass matrix (13) for $W^3, F^3, Y$ and the diagonal vectors (14) in [4] are also suitable for our model

\[
\begin{pmatrix}
  g^2_w B & g_w g_F C & g_w g_Y B \\
  g_w g_F C & g^2_F D & g_Y g_F C \\
  g_w g_Y B & g_Y g_F C & g^2_Y B
\end{pmatrix}
\] (13)

The diagonal neutral vectors are

\[
\begin{align*}
A^\mu &= \frac{1}{N_A} (g_Y W^\mu_3 - g_w Y^\mu) \\
Z^\mu &= \frac{1}{N_Z} (g_Y Y^\mu + g_w W^\mu_3 + g_F Y_1 F^\mu_3) \\
G^\mu &= \frac{1}{N_G} (g_Y Y^\mu + g_w W^\mu_3 + g_F Y_2 F^\mu_3)
\end{align*}
\] (14)

For conveniently comparing with the Weinberg angle in SM, we assume

\[
\begin{align*}
g_F &= g \cos \theta \\
g_Y &= g \sin \theta \sin \phi \\
g_w &= g \sin \theta \cos \phi
\end{align*}
\] (15)

The interaction Lagrangian can then be written as
\[ \mathcal{L}_{int} = \frac{\sqrt{2}}{2} g_w [W_\alpha^- (\bar{e}_L \gamma^\alpha v_e + \bar{\mu}_L \gamma^\alpha v_\mu) + W_\alpha^+ (\bar{e}_e \gamma^\alpha e_L + \bar{\mu}_\mu \gamma^\alpha \mu_L)] \\
- \frac{\sqrt{2}}{2} G_F [F^- \bar{e}_e \gamma^\alpha v_e + F^+ \bar{\mu}_\mu \gamma^\alpha v_\mu] \\
+ g \sin \theta \sin \phi \cos \phi A_\alpha (\bar{e} \gamma^\alpha e + \bar{\mu} \gamma^\alpha \mu) \\
+ \frac{y_2}{y_2 - y_1} (\cos^2 \phi - \sin^2 \phi) (\bar{e}_L \gamma^\alpha e_L + \bar{\mu}_L \gamma^\alpha \mu_L) \cdot N_Z Z_\alpha \\
+ \frac{1}{2} \left( \frac{1 - y_2}{y_2 - y_1} \bar{e}_e \gamma^\alpha v_e - \frac{1 + y_2}{y_2 - y_1} \bar{\mu}_\mu \gamma^\alpha v_\mu \right) \cdot N_Z Z_\alpha \\
+ \frac{y_1}{y_2 - y_1} \sin^2 \phi (\bar{R}_e \gamma^\alpha R_e + \bar{R}_\mu \gamma^\alpha R_\mu) \cdot N_G G_\alpha \\
+ \frac{1}{2} \left( \frac{y_1}{y_2 - y_1} \sin^2 \phi - \frac{1}{y_2 - y_1} \right) \bar{e}_e \gamma^\alpha v_e \cdot N_G G_\alpha \\
+ \frac{1}{2} \left( \frac{y_1}{y_2 - y_1} \sin^2 \phi + \frac{1}{y_2 - y_1} \right) \bar{\mu}_\mu \gamma^\alpha v_\mu \cdot N_G G_\alpha \tag{16} \]

The coupling constants in (16) compared with the ones in SM will yield

\[ e = g \sin \theta \sin \phi \cos \phi \]

\[ \frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_w^2} \sin^2 \theta \cos^2 \phi \tag{17} \]

Comparing the third formula in (15) and the second formula in (17), we can deduce that \( g_w \) in our model is just the one in SM. Combined with the first formula in (17), we can deduce that \( \phi \) is just the Weinberg angle \( \theta_W \) in SM. The maximal mixture between neutrinos implies the coupling constant \( g_F \) for interactions between different neutrinos is similar to \( g_Z \) in SM for the same flavors. Then we assume the magnitude of \( g_F \) is the same as \( g_Z \) in SM.

**The flavor conversion probability per unit time**

**The cross section \( \sigma \) and the scattering probability per unit time at the start of neutrino flight**

We assume neutrinos collide in beam due to their ‘Brownian movement’ with their moments of all possible directions. According to our model, neutrino flavor conversions are caused by interactions intermediated by \( F_\mu \). At the start of neutrino’s flight, all neutrino-neutrino interactions can be shown in one Feynman diagram as (18)
$\nu_\alpha \rightarrow k \rightarrow N$ \hspace{.5cm} $\nu_\beta, \gamma = \epsilon, \mu, \tau$ \hspace{.5cm} (18)

$v_\alpha$ is the initial neutrino and $N$ is vector boson which intermediates neutrino-neutrino interactions such as $Z, F^{\pm}$. We assume these vector bosons have masses similar to $Z^0$ in SM.

For small mass of neutrino, our following calculation will ignore neutrino mass. Thus, the scattering amplitude of (18) is

$$i M = -\frac{ig_N^2}{M_N^2} [\bar{u}_{\nu_\beta}(k)\gamma^\mu(1-\gamma_5)u_{\nu_\alpha}(p)] [\bar{u}_{\nu_\gamma}(k')\gamma^\mu(1-\gamma_5)u_{\nu_\alpha}(p')]$$ \hspace{.5cm} (19)

$m_{\nu_\epsilon} \sim m_{\nu_\mu} \sim m_{\nu_\tau} \sim 0$ \hspace{.5cm} (20)

As is mentioned, the magnitude of $g_F \sim g_Z (in SM)$ and then

$$\sigma_{CM} = \frac{|M|^2}{16\pi E_{cm}^2} = \frac{4^3 \times g_F^4}{16\pi E_{cm}^2 M_N^4} (k \cdot k')(p \cdot p') = \frac{4 \times g_F^4}{\pi E_{cm}^2 M_N^4} (k \cdot k')(p \cdot p')$$ \hspace{.5cm} (21)

Not that (21) is obtained with the convention that the impact parameter $b \sim 0$.

The neutrino conversion probability generally is obtained on axis, so (21) makes sense.

In center of mass frame, we have

$$k \cdot k' = \frac{1}{2} (2k \cdot k') = \frac{1}{2} [(k + k')^2 - (m_\beta^2 + m_\gamma^2)]$$

$$p \cdot p' = \frac{1}{2} (2p \cdot p') = \frac{1}{2} [(p + p')^2 - 2m_\alpha^2]$$

And we get scattering cross section in center of mass frame as below

$$\sigma_{CM} = \frac{4 \times g_F^4}{\pi E_{cm}^2 M_N^4} (k \cdot k')(p \cdot p') = \frac{g_F^4 \times [E_{cm}^2 - (m_\beta^2 + m_\gamma^2)] \times [E_{cm}^2 - 2m_\alpha^2]}{\pi E_{cm}^2 M_N^4} \approx \frac{g_F^4 \times E_{cm}^2}{\pi M_N^4}$$ \hspace{.5cm} (22)

Let

$$s = (p + p')^2 = (k + k')^2 = E_{cm}^2$$ \hspace{.5cm} (23)

We can obtain the expression about $s$ of section $\sigma$
Taking the rest frame of the initial neutrino $v_\alpha(p)$ before interaction, we have

$$s = 2p \cdot p' + 2m_{v_\alpha}^2 \approx 2p \cdot p' \approx 2m_{v_\alpha}E_{v_\alpha}$$  \hspace{1cm} (25)$$

Then we have scattering cross section as below

$$\sigma \approx \frac{8^4}{\pi M_z^4} \times s \approx \frac{8^4 \times 2m_{v_\alpha}E_{v_\alpha}}{\pi M_z^4}$$  \hspace{1cm} (26)$$

After some calculation of natural units conversion (which we will show in Appendix B), we get

$$\sigma \approx 10^{-19} \left( \text{GeV} \right)^{-2} = \frac{1}{4\pi} \times 10^{-47} \text{cm}^2$$  \hspace{1cm} (27)$$

Here we assume $m_{v_\alpha} = m_{v_e} \sim 1\text{eV}$ and the energy of neutrino $E = 1\text{GeV}$. Generally, the masses of the former two neutrino families $v_e, v_\mu$ are considered far lighter than the third $v_\tau$. Thus when $v_\tau$ participates in interactions, its mass in (22) cannot be ignored and (27) is not suitable for these interactions and therefore the cross section of these interactions should be obtained from (22) directly. In this paper, we only want to know whether our predicted conversion probability is consistent with experiments. For simplicity and without loss of generality, we only need to evaluate the order of predicting $P(v_e \rightarrow v_\mu)$.

Following, we will show how to evaluate conversion probability by cross section $\sigma$. Consider the definition of $\sigma$

$$\sigma = \frac{N}{n_B \cdot N_A} = \frac{1}{n_B \cdot N_A} \int d^2 b \: n_B P(b).$$  \hspace{1cm} (28)$$

In equation (28), $n_B$ denotes current density of incident particle B; $N_A$ denotes particle number of target particle A; N denotes all scattering events in unit time.

From (28), if beam section is unit area, the value of current density $n_B$ equals incident particle number in unit time. In neutrino beam, target particle is also particles in beam of the same section and the number of target particle $N_A = n_B$. Thus, according to (28), the value of $\sigma$ equals the average scattering probability per neutrino in unit time. That means the value of $\sigma$ equals the value of average conversion probability per unit time of one initial neutrino. Because the cross section formula (21) is established with the convention $b \sim 0$, we will choose the unit area 1cm$^2$. Then
the magnitude of neutrino-neutrino scattering probability per unit time is

\[ P(v_\alpha + v_\alpha \rightarrow v_\beta + v_\gamma) = \frac{1}{4\pi} \times 10^{-47} \times 2 \]  

(29)

The factor 2 is because 2 initial neutrinos.

Thus our predicting neutrino flavor conversion probability per unit time is

\[ P(v_\alpha \rightarrow v_\beta) \sim O(10^{-47}) \]  

(30)

The actual conversion probability per unit time at the start of flight

We can obtain the conversion probability per time (32) by taking a derivative of the empirical formula (31) with respect to t:

\[ P(v_\alpha \rightarrow v_\beta) = \sin^2(2\theta_{ij})\sin^2\left(\frac{t}{L_0}\right) \quad L_0 = \frac{4E}{\Delta m^2_{ij}} \]  

(31)

\[ \frac{dP(v_\alpha \rightarrow v_\beta)}{dt} = 2\sin^2(2\theta_{ij})\sin\left(\frac{t}{L_0}\right)\cos\left(\frac{t}{L_0}\right) \frac{1}{L_0} \approx 2\sin^2(2\theta_{ij}) \frac{t}{L_0^2} \]  

(32)

The last term of (32) is because \( t \ll L_0 \) at the start of flight.

In the first second of flight,

\[ \frac{dP(v_\alpha \rightarrow v_\beta)}{dt} \approx 2\sin^2(2\theta_{ij}) \frac{t}{L_0^2} = 2\sin^2(2\theta_{ij}) \frac{1}{L_0^2} = 2\sin^2(2\theta_{ij}) \left(\frac{\Delta m^2_{ij}}{4E}\right)^2 \]  

(33)

To estimate the magnitude of (33), we take a special example as the conversion of \( v_e \rightarrow v_\mu \). According to experimental data, we take values of (34) into (33):

\[ E \approx 1GeV \]
\[ \Delta m^2_{12} \approx 8 \times 10^{-5}(eV)^2 = 8 \times 10^{-23}(GeV)^2 \]
\[ \sin^2(2\theta_{12}) = 0.86 \]  

(34)

So, the conversion probability per unit time at the start of neutrino flight is

\[ P(v_e \rightarrow v_\mu) \approx 2\sin^2(2\theta_{12})\left(\frac{8 \times 10^{-23}}{4}\right)^2 \approx 6 \times 10^{-46} \]  

(35)

Apparently, our prediction (30) is very close to the actual value(35).
Neutrino oscillation

According to our model, neutrino-neutrino interaction is

$$L_{\text{int}} = -ig_Z' \sum_{\alpha} \bar{v}_\alpha \gamma^\mu (1 - \gamma_5) v_\alpha Z_\mu - ig_F' \sum_{\alpha \neq \beta} \bar{v}_\beta \gamma^\mu (1 - \gamma_5) v_\alpha F^\pm_\mu$$

(36)

As is mentioned before, we assume

$$g'_Z \sim g_Z \sim g_F, \quad M_Z \sim M_F,$$

(37)

then the whole approximate effective Hamiltonian can be written in one term as follows

$$H_{\text{int}} = \frac{G'_F}{\sqrt{2}} \int dx^3 \left( \sum_i \bar{v}'_i L \gamma^\mu v_i L \right) \left( \sum_j \bar{v}'_j L \gamma^\mu v_j L \right)$$

(38)

Here $v'$ is not a pure state and is some mixture of several states. While $v_i L, v_j L$ are initial states. Thus the neutrino mixture is just the macro result of the sum of all interactions. And we also can conclude that the mixture of neutrino flavors will not be constant until all neutrino-neutrino interactions arrive at balance.

We can see that there is an equilibrium state for these interactions. There must have some oscillations before neutrino probability distribution arrive at equilibrium distribution.

Conclusions

For interactions responsible for neutrino flavor conversions, we constructed the extension of SM including a horizontal flavor symmetry i.e. $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$. The horizontal symmetry $SU(2)_N$ introduces interactions between different neutrino flavors and leads to neutrino flavor conversions.

To testify our idea, we have evaluated the flavor conversion probability induced by new interactions by utilizing the definition of cross section. In this calculation, we assume neutrinos collide due to 'Brown movement' and the direction of individual neutrino instantaneous moment in beam is not definite and generally not consistent to the direction of neutrino flux. Fortunately, the prediction is consistent with experimental data. And we also demonstrate neutrino oscillation as phenomenon before all neutrino-neutrino interactions arriving at balance.

Appendix

A. The invariance of the fermion dynamical Lagrangian under the group $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$

The invariance of $L$ transforming under the group $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$ can be equivalent to invariance under group $SU(2)_L, U(1)_Y$ and $SU(2)_N$ in the
We will prove it term by term.

$$L \rightarrow e^{-i \frac{1}{2} \alpha^a(x) \tau^a_w} L$$
$$L \rightarrow e^{-i \frac{1}{2} \beta(x) \hat{Y}} L \quad (A.1)$$
$$L \rightarrow L e^{i \frac{1}{2} \rho^b(x) \tau^b_F}$$

\(\alpha, \beta, \rho\) is infinitesimal, \(\tau_w, \tau_F\) are generators associated with \(SU(2)_L\) and \(SU(2)_N\) respectively. \(\hat{Y}\) is the weak hypercharge operator related to \(U(1)_Y\). We introduce gauge fields \(W^i_a, Y_\mu, F^i_a\) to relate groups \(SU(2)_L, U(1)_Y\) and \(SU(2)_N\) respectively.

The covariant derivatives related with transformation under \(SU(2)_L, U(1)_Y\) and \(SU(2)_N\) respectively are

$$D_\mu = \partial_\mu - i \frac{1}{2} \gamma^a \tau^a_w W^a_\mu$$
$$D_\mu = \partial_\mu - i \frac{1}{2} \gamma^Y Y_\mu \hat{Y} \quad (A.2)$$
$$D_\mu L = \partial_\mu L + i \frac{1}{2} \gamma^F \tau^a_F F^a_\mu$$

Define the gauge transformation of gauge fields \(W^i_a, Y_\mu, F^i_a\) as follows

$$i \gamma^a \tau^a W^a_\mu = i \gamma^a \tau^a W^a_\mu - i \partial_\mu a^b \tau^b - \frac{1}{2} \gamma^w [\tau^a W^a_\mu, \alpha^b \tau^b]$$
$$Y_\mu = Y_\mu - \frac{i}{2} \gamma^Y \partial_\mu \beta$$
$$i \gamma^F \tau^a F^a_\mu = i \gamma^F \tau^a F^a_\mu - i \partial_\mu \rho^b \tau^b - \frac{1}{2} \gamma^F [\tau^a F^a_\mu, \rho^b \tau^b] \quad (A.3)$$

The fermion dynamical Lagrangian invariant under the group \(SU(2)_L \otimes U(1)_Y \otimes SU(2)_N\) is

$$\mathcal{L}_f = i \bar{R}_c \gamma^a (\partial_\alpha + i \gamma Y \alpha) R_c + i \bar{R}_\mu \gamma^a (\partial_\alpha + i \gamma Y \alpha) R_\mu$$
$$+ i \mathcal{I}[\gamma^a (\partial_\alpha + i \gamma Y \alpha - i \frac{1}{2} \gamma^w \tau^a_w W^a_\mu) L]$$
$$+ i \mathcal{I}[\gamma^a i \frac{1}{2} \gamma^F (1 + \tau_3) L \tau^a_F F^a_\mu] \quad (A.4)$$

In (A.4), for left-handed lepton, the value of the hypercharge \(Y\) is -1 and for right-handed lepton \(R\), the value of the hypercharge \(Y\) is -2.

With the transformation of these gauge fields (42), the fermion dynamical Lagrangian (A.4) will be invariant under the group of \(SU(2)_L \otimes U(1)_Y \otimes SU(2)_N\). We will prove it term by term.

**Proof:**

1. The first two terms in (A.4)

$$R_c \rightarrow e^{-i \frac{1}{2} \beta(x) \hat{Y}} R_c \rightarrow (1 - i \frac{1}{2} \beta(x) \hat{Y}) R_c = (1 + i \beta(x)) R_c$$

The covariant derivative is

$$D_\mu = \partial_\mu + i \gamma^Y Y_\mu$$
and the transformation of gauge field

\[ Y_\mu \to Y_\mu - \frac{1}{g_Y} \partial_\mu \beta \]

We only need to prove

\[ D_\mu (1 + i \beta(x)) R_e \to (1 + i \beta(x)) D_\mu R_e \]

i.e.

\[ (\partial_\mu + ig_Y Y_\mu) [(1 + i \beta(x)) R_e] \to (1 + i \beta(x)) (\partial_\mu + ig_Y Y_\mu) R_e \quad \text{(A.5)} \]

Left of (A.5) =

\[
i \partial_\mu \beta R_e + (1 + i \beta) \partial_\mu R_e + ig_Y Y_\mu R_e - g_Y Y_\mu \beta R_e \\
\to i \partial_\mu \beta R_e + (1 + i \beta) \partial_\mu R_e + ig_Y (Y_\mu - \frac{1}{g_Y} \partial_\mu \beta) R_e - g_Y (Y_\mu - \frac{1}{g_Y} \partial_\mu \beta) \beta R_e \\
\to \partial_\mu R_e + i \partial_\mu \beta R_e + ig_Y Y_\mu R_e - i \partial_\mu \beta R_e - g_Y Y_\mu \beta R_e + \partial_\mu \beta \beta R_e \\
\to \partial_\mu R_e + i \beta \partial_\mu R_e + ig_Y Y_\mu R_e - g_Y Y_\mu \beta R_e + \partial_\mu \beta \beta R_e
\]

Right =

\[ (1 + i \beta(x)) (\partial_\mu + ig_Y Y_\mu) R_e \]

\[ \to \partial_\mu R_e + i \beta \partial_\mu R_e + ig_Y Y_\mu R_e - g_Y Y_\mu \beta R_e \]

Omitting the higher order of \( O(g\beta) \), we can obtain

\[ (\partial_\mu + ig_Y Y_\mu) [(1 + i \beta(x)) R_e] \to (1 + i \beta(x)) (\partial_\mu + ig_Y Y_\mu) R_e \]

2. The third and fourth term

We will prove the invariance of these two terms by prove invariance under group \( SU(2)_L \), \( U(1)_Y \) and \( SU(2)_N \) respectively.

The transformations of the gauge fields are

\[
\tau^a W^a_\mu \to \tau^a W^a_\mu - \frac{1}{g_W} \partial_\mu \alpha^b \tau^b + i \frac{1}{2} [\tau^a W^a_\mu, \alpha^b \tau^b] \\
Y_\mu \to Y_\mu - \frac{1}{g_Y} \partial_\mu \beta \\
\tau^a F^a_\mu \to \tau^a F^a_\mu - \frac{1}{g_F} \partial_\mu \rho^b \tau^b + i \frac{1}{2} [\tau^a F^a_\mu, \rho^b \tau^b]
\]

(1) (A.4) is invariant under \( U(1)_Y \)
\[ L \rightarrow [1 + i \frac{1}{2} \beta(x)] L \]
\[ D_\mu = \partial_\mu + i \frac{1}{2} g Y_\mu \]

Which we need to prove is
\[ L \rightarrow [1 + i \frac{1}{2} \beta(x)] L \]
\[ D_\mu [1 + i \frac{1}{2} \beta(x)] L \rightarrow [1 + i \frac{1}{2} \beta(x)] D_\mu L, \]

that is
\[ (\partial_\mu + i \frac{1}{2} g Y_\mu)[1 + i \frac{1}{2} \beta(x)] L \rightarrow [1 + i \frac{1}{2} \beta(x)](\partial_\mu + i \frac{1}{2} g Y_\mu) L \]  \hspace{1cm} (A.6)

with
\[ Y_\mu \rightarrow Y_\mu - \frac{1}{g} \partial_\mu \beta \]

Left of (A.6) =
\[ (\partial_\mu + i \frac{1}{2} g Y_\mu)[1 + i \frac{1}{2} \beta(x)] L \]
\[ = \frac{1}{2} \partial_\mu \beta(x) L + [1 + \frac{1}{2} \beta(x)] \partial_\mu L + i \frac{1}{2} g Y_\mu L - \frac{1}{4} g Y \beta \]
\[ \approx [1 + i \frac{1}{2} \beta(x)] \partial_\mu L + i \frac{1}{2} g Y_\mu L \]

The right of (A.6) is
\[ [1 + i \frac{1}{2} \beta(x)](\partial_\mu + i \frac{1}{2} g Y_\mu) L, \]
\[ = [1 + i \frac{1}{2} \beta(x)] \partial_\mu L + i \frac{1}{2} g Y_\mu L - \frac{1}{4} g Y \beta \]
\[ \approx [1 + i \frac{1}{2} \beta(x)] \partial_\mu L + i \frac{1}{2} g Y_\mu L \]

Thus (A.6) can be established.

(2) (A.4) is invariant under \( SU(2)_L \)
\[ L \rightarrow [1 - i \frac{1}{2} \alpha^a(x) \tau^a] L \]
\[ D_\mu = \partial_\mu - i \frac{1}{2} g \omega^a \tau^a W^a_\mu \]
Which we need to prove is

\[ D_{\mu}[1 - i \frac{1}{2} a^a(x) \tau^a] L = [1 - i \frac{1}{2} a^a(x) \tau^a] D_{\mu} L, \]

that is

\[ (\partial_{\mu} - i \frac{1}{2} g_w \tau^a W_{\mu}^a) [1 - i \frac{1}{2} a^a(x) \tau^a] L = [1 - i \frac{1}{2} a^a(x) \tau^a] (\partial_{\mu} - i \frac{1}{2} g_w \tau^a W_{\mu}^a) L \quad (A.7), \]

with

\[ \tau^a W_{\mu}^a \rightarrow \tau^a W_{\mu}^a - \frac{1}{g_w} \partial_{\mu} \alpha^b \tau^b + i \frac{1}{2} [\tau^a W_{\mu}^a, \alpha^b \tau^b] \]

The left of (A.7) is

\[
\begin{align*}
(\partial_{\mu} - i \frac{1}{2} g_w \tau^a W_{\mu}^a) [1 - i \frac{1}{2} a^a(x) \tau^a] L \\
= -i \frac{1}{2} \partial_{\mu} a^a(x) \tau^a L + [1 - i \frac{1}{2} a^a(x) \tau^a] \partial_{\mu} L - i \frac{1}{2} g_w \tau^a W_{\mu}^a L - \frac{1}{4} g_w \tau^a W_{\mu}^a a^a(x) \tau^a L \\
\rightarrow -i \frac{1}{2} \partial_{\mu} a^a(x) \tau^a L + [1 - i \frac{1}{2} a^a(x) \tau^a] \partial_{\mu} L - i \frac{1}{2} g_w \{ \tau^a W_{\mu}^a - \frac{1}{g_w} \partial_{\mu} \alpha^b \tau^b + i \frac{1}{2} [\tau^a W_{\mu}^a, \alpha^b \tau^b] \} L \\
- \frac{1}{4} g_w \{ \tau^a W_{\mu}^a - \frac{1}{g_w} \partial_{\mu} \alpha^b \tau^b + i \frac{1}{2} [\tau^a W_{\mu}^a, \alpha^b \tau^b] \} a^c(x) \tau^c L \\
= -i \frac{1}{2} \partial_{\mu} a^a(x) \tau^a L + [1 - i \frac{1}{2} a^a(x) \tau^a] \partial_{\mu} L - i \frac{1}{2} g_w \tau^a W_{\mu}^a L + i \frac{1}{2} \partial_{\mu} \alpha^b \tau^b L \\
+ \frac{1}{4} g_w [\tau^a W_{\mu}^a, \alpha^b \tau^b] L - \frac{1}{4} g_w \tau^a W_{\mu}^a \alpha^b \tau^b L \\
= [1 - i \frac{1}{2} a^a(x) \tau^a] \partial_{\mu} L - i \frac{1}{2} g_w \tau^a W_{\mu}^a L + \frac{1}{4} g_w [\tau^a W_{\mu}^a, \alpha^b \tau^b] L - \frac{1}{4} g_w \tau^a W_{\mu}^a \alpha^b \tau^b L \\
= [1 - i \frac{1}{2} a^a(x) \tau^a] \partial_{\mu} L - i \frac{1}{2} g_w \tau^a W_{\mu}^a L - \frac{1}{4} g_w \alpha^b \tau^b \tau^a W_{\mu}^a L.
\end{align*}
\]

Here, we omitting the higher orders.

The right of (A.7)=

\[
[1 - i \frac{1}{2} a^a(x) \tau^a] (\partial_{\mu} - i \frac{1}{2} g_w \tau^a W_{\mu}^a) L \\
= [1 - i \frac{1}{2} a^a(x) \tau^a] \partial_{\mu} L - i \frac{1}{2} g_w \tau^a W_{\mu}^a L - \frac{1}{4} g_w \alpha^b \tau^b \tau^a W_{\mu}^a L.
\]

Thus (A.7) can be established.

(3) (A.4) is invariant under the horizontal group $SU(2)_N$

\[
L \rightarrow L' = L (1 + i \frac{1}{2} \rho^b(x) \tau^b) \\
D_{\mu} L = \partial_{\mu} L + i \frac{1}{2} g_F L \tau^a F_{\mu}^a
\]
Which we need to prove is

\[ D_\mu [L(1 + i \frac{1}{2} \rho^b(x) \tau^b)] \rightarrow (D_\mu L)(1 + i \frac{1}{2} \rho^b(x) \tau^b), \]

that is

\[
\partial_\mu [L(1 + i \frac{1}{2} \rho^b(x) \tau^b)] + i \frac{1}{2} GF L(1 + i \frac{1}{2} \rho^b(x) \tau^b) \tau^a F^a_\mu \\
\rightarrow (\partial_\mu L + i \frac{1}{2} GF L \tau^a F^a_\mu)(1 + i \frac{1}{2} \rho^b(x) \tau^b)
\]

(A.8)

with

\[
\tau^a F^a_\mu \rightarrow \tau^a F^a_\mu - \frac{1}{8F} \partial_\mu \rho^b \tau^b + i \frac{1}{2} \left[ \tau^a F^a_\mu, \rho^b \tau^b \right]
\]

The left of (A.8) is

\[
\partial_\mu [L(1 + i \frac{1}{2} \rho^b(x) \tau^b)] + i \frac{1}{2} GF L(1 + i \frac{1}{2} \rho^b(x) \tau^b) \tau^a F^a_\mu \\
= \partial_\mu L(1 + i \frac{1}{2} \rho^b(x) \tau^b) + Li \frac{1}{2} \partial_\mu \rho^b(x) \tau^b + i \frac{1}{2} GF L \tau^a F^a_\mu - \frac{1}{4} GF L \rho^b(x) \tau^b \tau^a F^a_\mu \\
\rightarrow \partial_\mu L(1 + i \frac{1}{2} \rho^b(x) \tau^b) + Li \frac{1}{2} \partial_\mu \rho^b(x) \tau^b + i \frac{1}{2} GF L \{ \tau^a F^a_\mu - \frac{1}{8F} \partial_\mu \rho^b \tau^b + i \frac{1}{2} \left[ \tau^a F^a_\mu, \rho^b \tau^b \right] \} \\
- \frac{1}{4} GF L \rho^c(x) \tau^c \{ \tau^a F^a_\mu - \frac{1}{8F} \partial_\mu \rho^b \tau^b + i \frac{1}{2} \left[ \tau^a F^a_\mu, \rho^b \tau^b \right] \} \\
\approx \partial_\mu L(1 + i \frac{1}{2} \rho^b(x) \tau^b) + Li \frac{1}{2} \partial_\mu \rho^b(x) \tau^b + i \frac{1}{2} GF L \tau^a F^a_\mu - i \frac{1}{2} \partial_\mu \rho^b \tau^b \\
- \frac{1}{4} GF L \tau^a F^a_\mu \rho^b \tau^b \\
= \partial_\mu L(1 + i \frac{1}{2} \rho^b(x) \tau^b) + i \frac{1}{2} GF L \tau^a F^a_\mu - \frac{1}{4} GF L \tau^a F^a_\mu \rho^b \tau^b
\]

The right of (A.8) is

\[ (\partial_\mu L + i \frac{1}{2} GF L \tau^a F^a_\mu)(1 + i \frac{1}{2} \rho^b(x) \tau^b) \]

\[ = \partial_\mu L(1 + i \frac{1}{2} \rho^b(x) \tau^b) + i \frac{1}{2} GF L \tau^a F^a_\mu - \frac{1}{4} GF L \tau^a F^a_\mu \rho^b \tau^b \]

Thus (A.8) can be established.

So, (A.4) is invariant under $SU(2)_L \otimes U(1)_Y \otimes SU(2)_N$.

**B. The cross section of neutrino-neutrino interaction $\sigma$**

We assign neutrino energy $E_\nu = 1$GeV, and some known parameters as below

\[
m_{\nu} \approx 1eV = 10^{-9}$GeV
\]
\[
M_Z \approx 90$GeV
\]
\[
S_Z^2 = \frac{M_Z}{4\sqrt{2}}GF
\]
\[
G_F = 1 \times 10^{-5}/(\text{GeV})^2
\]
Taking above values into the cross section formula, we will get

\[
\sigma \approx \frac{g^4 \times 2m_{\nu_e}E_{\nu_e}}{\pi M^4_z} = \frac{M^4_z}{32} \times \frac{G_F^2 \times 2m_{\nu_e}E_{\nu_e}}{\pi M^4_z} = \frac{10^{-19}}{16\pi} (GeV)^{-2} \tag{B.1}
\]

Using the conversion of natural units, we can get

\[
1 cm = 5 \times 10^{13} GeV^{-1} \\
\rightarrow GeV^{-1} = 0.2 \times 10^{-13} cm \\
\rightarrow GeV^{-2} = 0.04 \times 10^{-26} cm^2
\]

Thus

\[
\sigma \approx \frac{10^{-19}}{16\pi} (GeV)^{-2} = \frac{1}{4\pi} \times 10^{-47} cm^2
\]

Because the cross section formula (B.1) is obtained with the convention that the impact parameter \( b \sim 0 \), we adopt the unit \( cm^2 \).

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