Monolayer Sc$_2$CF$_2$ as a Potential Selective and Sensitive NO$_2$ Sensor: Insight from First-Principles Calculations
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ABSTRACT: Two-dimensional materials with excellent surface–volume ratios and massive reaction sites recently have been receiving attention for gas sensing. With first-principles calculations, we explored the performance of monolayer Sc$_2$CF$_2$ as a gas sensor. We investigated how molecule adsorption affects its electronic structure and optical properties. It is found that a large charge transfer quantity happens between Sc$_2$CF$_2$ and NO$_2$, which results from the fact that the lowest unoccupied molecular orbital (LUMO) of NO$_2$ is below the valence band maximum (VBM) of Sc$_2$CF$_2$. Moreover, the MD simulation shows that NO$_2$ can adsorb on the Sc$_2$CF$_2$ surface stably at room temperature. We explored the effect of biaxial strain on the adsorption energy and charge transfer quantity of each system, and the results show that the biaxial strain can enhance both the adsorption energy and charge transfer quantity of the NO$_2$ system and thus can improve the sensitivity of Sc$_2$CF$_2$ in detecting the NO$_2$ molecule. Furthermore, we investigated the adsorption behavior and charge transfer of polar polyatomic molecules at the Sc$_2$CF$_2$ surface with h-BN as a substrate, and the results demonstrate that the h-BN substrate can hardly modify the main results. Our result predicts that Sc$_2$CF$_2$ can be a promising selective and sensitive sensor to detect the NO$_2$ molecule, and could also give a theoretical guide for other terminated MXenes used for gas sensors or detectors.

1. INTRODUCTION
Gas sensors and detectors play an important role in environmental protection and daily health safety applications. Because of the advantages of low costs, metal oxides have been explored widely; however, they require a high operating temperature. An excellent surface–volume ratio and massive reaction sites, a large number of 2D materials including graphene, h-BN, phosphorene, transition-metal chalcogenides, and layered group III–VI semiconductors have been reported with superior performance for detecting gas. Charge transfer happens when molecules are adsorbed on 2D materials, and the charge transfer between molecules and 2D materials will modify the electronic structures and/or resistance of sensing materials by varying degrees. A copious amount of charge transfer is required for high selectivity and sensitivity 2D sensing materials.

MXenes, including traditional metal carbides, nitrides, and carbon-nitrides with possible stoichiometric ratios of 2:1, 3:2, and 4:3, are a large family of 2D materials, which are usually synthesized through chemical exfoliation from the MAX phase by etching the ‘A’ atoms in acid solutions. Meanwhile, the intrinsic metallic surface of MXenes is usually passivated or terminated by a function group, and there consist four passivation models depending on function group adsorption sites on both sides. With a natural layered structure, MXenes have been explored as high-performance electrode materials in lithium-ion batteries. For instance, bare Ti$_2$C can show excellent theoretical capacity and low Li-ion diffusing barrier, and the rate performance of terminated Ti$_2$C can be greatly enhanced after intercalation. MXenes also find potential applications in other energy storage aspects, like supercapacitors, Li-S batteries, and multivalent-ion batteries. Owing to their intrinsic metallic property, MXenes flakes, which have low average resistivity and high electron mobility, also have been reported to be used as a conductive channel in field-effect transistors, which may construct an Ohmic contact. Besides, low-thickness Ti$_2$CX films with a transmittance of about 90% have been explored as transparent conductive electrodes. It has also been reported that active sites for water redox exist in MXenes, which also show suitable Gibbs free energy for hydrogen adsorption and are electro-catalyst or photo-catalyst candidates for water splitting. MXenes have found promising application in various fields.

Besides, MXenes also have been explored as sensing materials for sensors or detectors. In terms of theoretical calculations, MXenes have been widely explored for gas sensing. Density functional theory (DFT) calculations demonstrated that bare MXenes are reactive, where molecules’...
dissociative adsorption happens at most surfaces, and the selectivity is enhanced upon surface functionalization.\textsuperscript{48} DFT calculations along with non-equilibrium Green’s function (NEGF) method simulation predicted that O-terminated Ti$_2$C can detect NH$_3$ sensitively with a strong binding energy. A dramatic change of the $I$–$V$ curve after NH$_3$ adsorption determines the selectivity of O-terminated Ti$_2$C.\textsuperscript{49} In addition, applied strain can increase the binding energy and improve the sensitivity, while releasing the strain or injecting an electron into MXenes will weaken the strength and achieve reusability of sensing materials.\textsuperscript{49,50} Similarly, it has been reported that SO$_2$ can be detected by Sc$_2$CO$_2$ with suitable adsorption energy. Moreover, biaxial strain and external vertical electric field can help SO$_2$ desorption from Sc$_2$CO$_2$.\textsuperscript{51} In terms of experiments, Wang et al. have reported that stacked Ti$_3$C$_2$T$_x$ molecular sieving membranes exhibit excellent H$_2$ permeability, which mainly benefit from the sub-nanometer interlayer spacing between the neighboring MXene nanosheets.\textsuperscript{52} Stanciu et al. designed a nanohybrid structure with MXene (Ti$_3$C$_2$T$_x$) and transition metal dichalcogenide (WSe$_2$), which can adsorb oxygen species (O$_2$ and O$^-$) in air and exhibits good flexibility for O-containing volatile organic compounds with low noise and fast response/recovery speed.\textsuperscript{53}

To our best knowledge, MXenes are potential candidates for gas sensors. Here, we performed DFT calculations to explore the adsorption of 11 common molecules on F-terminated Sc$_2$C (Sc$_2$CF$_2$) and their charge transfer, which have not been studied yet. We investigated how molecule adsorption affects their electronic structure and optical properties. It is found that the largest charge transfer quantity happens between Sc$_2$CF$_2$ and NO$_2$, which is because the lowest unoccupied molecular orbital (LUMO) of NO$_2$ is below the valence band maximum (VBM) of Sc$_2$CF$_2$, although the adsorption energy is not very large. The adsorption stability and sensitivity of Sc$_2$CF$_2$ sensing NO$_2$ are demonstrated through molecular dynamics simulation with an NVT ensemble, and applying biaxial strain is claimed to be a promising method for enhancing the sensitivity and selectivity. Besides, h-BN as a substrate can hardly modify the adsorption energy and charge transfer at the polar polyatomic molecules’ adsorption system and is a promising substrate. These results indicate that F-terminated Sc$_2$C is a good candidate for a NO$_2$ gas sensor, and give a theoretical guide for other terminated MXenes used for gas sensors or detectors.

### 2. RESULTS AND DISCUSSION

To simulate the Sc$_2$CF$_2$ monolayer exposed in air, we place 11 common molecules on the surface of Sc$_2$CF$_2$, and several adsorption configurations, i.e., top, bridge, and hollow sites of the three top F atoms, as well as the vertical and parallel orientation of each molecule, are considered. To characterize the interaction strength between the molecule and the 2D sheet, we define the adsorption energy ($E_{ad}$) as $E_{ad} = E_{tot} - E_{2D} - E_{gas}$, where $E_{tot}$ and $E_{2D}$ are the energies of the monolayer Sc$_2$CF$_2$ with and without adsorption of a gas molecule, respectively; $E_{gas}$ is the energy of a gas molecule in a cubic lattice with a lattice constant of 15 Å. The most stable configurations and $E_{ad}$ for each molecule system are shown and listed in Figure 1 and Table 1. The homo-nuclear diatomic molecules, like H$_2$, O$_2$, and N$_2$, will stay parallel to the substrate, and the weakest interaction with $E_{ad}$ ranging from $-0.07$ to $-0.09$ eV per molecule is found in these cases. The polar hetero-nuclear diatomic molecules, like NO and CO, also remain parallel with Sc$_2$CF$_2$, and the $E_{ad}$ is also small, however, it is a little stronger than that of the homo-nuclear diatomic molecule ($E_{ad} = -0.13$ and $-0.14$ eV for NO and CO, respectively), which is due to the strong electronegativity of the surface F atoms of the substrate. The CO$_2$ molecule has no polarity and contacts weakly with Sc$_2$CF$_2$. For polar polyatomic molecules ($H_2S$, $H_2O$, SO$_2$, NO$_2$, and NH$_3$), the adsorption is much stronger. The $E_{ad}$ of H$_2$O and H$_2$S is $-0.20$ and $-0.18$ eV per molecule, respectively, and the H in both cases will favor staying at the top of F, while the O in H$_2$O tends to stay at the top of Sc and S prefers the top of C. The NO$_2$ is more likely to stand vertically on the substrate, which mainly benefits from the sub-nanometer interlayer spacing between the neighboring MXene nanosheets.

![Figure 1](https://doi.org/10.1021/acsomega.1c06027)  
**Figure 1.** Most stable adsorption configurations for 11 molecules (light gray balls for Sc atoms, gray for C, light blue for F, white for H, dark blue for N, red for O, and yellow for S). The upper panels are the top view, and the lower panels are the side view of each molecule.

| Molecule | $E_{ad}$ (eV) | $D_z$ (Å) | $D$ (Å) | $CT$ ($10^{-3}$ e) | $E_v$ (eV) | $m_e^*$ ($m_0$) | $m_h^*$ ($m_0$) | $W_F$ (eV) |
|----------|---------------|----------|--------|-------------------|-----------|---------------|---------------|-----------|
| bare     | 1.02          | 0.98     | 2.32   | 4.71              |           |               |               |           |
| NH$_3$   | $-0.39$       | 1.54     | 2.53   | $-83.5$           | 0.99      | 0.71          | 2.03          | 4.57       |
| NO$_2$   | $-0.21$       | 2.33     | 2.93   | 100.3             | 0.99      | 0.90          | 2.13          | 5.17       |
| O$_2$    | $-0.21$       | 2.39     | 3.02   | 8.9               | 1.00      | 0.69          | 1.93          | 4.72       |
| H$_2$O   | $-0.20$       | 1.81     | 2.24   | $-19.7$           | 0.99      | 0.68          | 1.96          | 4.99       |
| H$_2$S   | $-0.18$       | 2.28     | 2.38   | $-4.7$            | 1.00      | 0.89          | 1.91          | 4.99       |
| CO$_2$   | $-0.16$       | 2.57     | 3.04   | 4.6               | 1.01      | 0.90          | 2.18          | 4.73       |
| CO       | $-0.14$       | 2.66     | 3.13   | $-2.7$            | 0.99      | 0.79          | 1.97          | 4.86       |
| NO       | $-0.13$       | 2.92     | 3.38   | 1.2               | 1.00      | 0.90          | 2.13          | 4.44       |
| O$_2$    | $-0.09$       | 2.64     | 3.21   | $-1.7$            | 0.99      | 0.81          | 1.92          | 4.99       |
| H$_2$    | $-0.07$       | 2.28     | 2.64   | $-3.9$            | 1.00      | 0.68          | 1.96          | 4.84       |
| N$_2$    | $-0.09$       | 3.19     | 3.58   | 0.5               | 1.00      | 0.93          | 2.08          | 4.75       |

$^*$Negative $CT$ values represent electron transfer from molecules to the Sc$_2$CF$_2$ sheet.
Figure 2. (a) Charge transfer from molecules to Sc$_2$CF$_2$ (calculated from Bader charge analysis). (b) Planar average charge density difference projected along the z direction. Spatial charge density differences of (c) NO$_2$ and (d) NH$_3$ adsorbed on Sc$_2$CF$_2$ with an isosurface of 0.05 e/Bohr$^3$. The purple and magenta regions indicate electron accumulation and depletion, respectively. The arrow specifies the direction of electron transfer.

SO$_2$ and NO$_2$, the O atoms favor to stay at the top or bridge of Sc. Actually, the O atom from different molecules shares the same character in all cases. The NH$_3$ shows the strongest adsorption behavior with $E_{ad}$ of $-0.39$ eV and vertical distance along the z direction of 1.54 Å. As shown in Figure 1, the NH$_3$ molecule favors to stack parallel on the 2D sheets, with the N atom staying at the top of Sc and three H atoms sharing symmetrically three hollow sites of neighboring F atoms. The coupling between the nanosheet and NH$_3$ will change the bond angles of H–N–H from 106.37 to 107.27°; meanwhile, the F atoms of nanosheets near the adsorption site move slightly away from the original position because the N partially shares the lone pair of electrons with Sc.

Once the molecule is adsorbed on the surface, the electronic structures of whole systems change and charge will redistribute between the molecule and nanosheet. Bader charge calculations are performed to determine the electron on each single atom so that we can see the charge transfer quantity induced by the gas molecule adsorption. As shown in Table 1 and Figure 2a, small amount of charge transfer (0.0005 $e$) has been found in diatomic molecules and non-polar cases (N$_2$, H$_2$, O$_2$, CO, NO, and CO$_2$), and more charge transfer can be found in polar polyatomic molecules. The H$_2$S, H$_2$O, and NH$_3$ act as electron donors and transfer 0.0047, 0.0197, and 0.0835 e to Sc$_2$CF$_2$, respectively, while SO$_2$ serves as an electron acceptor and obtains 0.0087 e from Sc$_2$CF$_2$. Especially, the largest charge transfer happens at the NO$_2$ system, and NO$_2$ takes 0.1003 electron from 2D materials, although the vertical distance between NO$_2$ and Sc$_2$CF$_2$ is larger than that of the NH$_3$ system. Furthermore, we calculated the charge density difference and projected it along the z direction as shown in Figure 2b. It can be seen that the polar polyatomic molecules’ adsorption systems show obvious oscillation at the interface region, especially the NH$_3$, NO$_2$, and H$_2$O systems, which is consistent with the result from Bader analysis. To show the details, the spatial charge density difference for NO$_2$ and NH$_3$ adsorption on the Sc$_2$CF$_2$ surface is plotted in Figure 2c,d, respectively. In the case of NO$_2$, all atoms at NO$_2$ and even the O atoms, which correspond with the three immediate neighboring positive peaks at about 13.5, 14, and 14.5 Å in Figure 2b, will accept electrons from the substrate and even from Sc atoms, which are located far away from the interface (see Figure 2c). The adsorption of NO$_2$ will induce a large interface dipole at the interface region, with electron accumulation at the NO$_2$ molecule side and electron depletion at the substrate side. As NH$_3$ attaches on Sc$_2$CF$_2$, electrons from the NH$_3$ molecule orbital, which are mainly from the three H atoms in Figure 2d, and correspond with the large valley near 14 Å in Figure 2b, transfer into Sc$_2$CF$_2$ and mainly into F atoms at the interface (see Figure 2d). Meanwhile, the adsorption of NH$_3$ will induce a surface dipole at Sc$_2$CF$_2$, with electron accumulation near the surface of F atoms and electron depletion near Sc atoms at the subsurface. In the case of H$_2$O, H atoms still prefer to donate electrons, which correspond the valley near 14.5 Å, and O atoms still like to accept electrons, which correspond the peak near 13 Å in Figure 2b. In total, few electrons will transfer from H$_2$O into the substrate, which also produces a small surface dipole at the surface of Sc$_2$CF$_2$. According to the transferred charge quantity, we can see the selectivity of Sc$_2$CF$_2$ to detect the NO$_2$ molecule.

To explain the physical causes of charge transfer, we calculated the band structures and density of states (DOS) of the adsorption system and extracted individual contributions of the molecule and substrate. Figure 3 shows the results for NH$_3$, SO$_2$, NO, and NO$_2$ adsorption systems. Several flat bands appear after adsorption, whose position has a close relationship with the alignment of frontier molecular orbitals of molecules and bands of pristine Sc$_2$CF$_2$. The isolated NH$_3$ molecule has a HOMO–LUMO gap of 5.4 eV from DFT calculation. Both the HOMO and LUMO of isolated NH$_3$ are far away from the VBM and conduction band minimum (CBM) of pristine Sc$_2$CF$_2$; thus, the flat bands of the NH$_3$ adsorption system stay away from the Fermi level. However, the shortest vertical distance between NH$_3$ and Sc$_2$CF$_2$ makes the strongest wavefunction mixing or electronic state coupling, which induces the large peaks at the energy range from $-2.5$ to $-0.3$ eV at the...
the SO2 adsorption system, the LUMO of the molecule is after adsorption, and the initial empty HOMO1 stays above the convenience. The two HOMO orbitals are mid-gap states large decrease of magnetic moment after adsorption (from 1 to figure 3c. The vertical distance is longer than that of the HOMO is located at the deep valance band region as seen located at the forbidden band region of Sc2CF2, while the HOMO is located at the deep valance band region as seen from figure 3c. The vertical distance is longer than that of the NH3 system; thus, there is less wave-function mixing, demonstrating much less charge transfer. For open-cell NO, the half-filled doubly degenerated antibonding HOMO (2π orbital with spin-up), together with its spin-split partner unoccupied LUMO (2π orbital with spin-down), becomes non-degenerated because of the symmetry breaking after contacting with the substrate. We named the non-degenerated HOMO orbital at the higher energy as HOMO1 and the other as HOMO2, and the non-degenerated LUMO at the lower energy level as LUMO1 and the other as LUMO2 for convenience. The two HOMO orbitals are mid-gap states after adsorption, and the initial empty HOMO1 stays above the Fermi level. As a polar hetero-nuclear diatomic molecule, the E_{ad} of NO is pretty small, which means weak wave-function mixing and is consistent with the small amount of charge transfer from the Bader calculation. It should be noticed that HOMO1 is still contributed from electrons of spin-up while HOMO2 is contributed from electrons of spin-down, which is totally different from the isolated NO molecule, indicating the large decrease of magnetic moment after adsorption (from 1 to 0.06 μB). The case of the NO2 adsorption system is very different from the three cases above, where the LUMO (6π, molecule orbital with spin-down) of NO2 is located below the VBM of Sc2CF2; thus, a large amount charge will transfer from the substrate to NO2 to fill the LUMO orbital of NO2. The valence band at some K points of the NO2 adsorption system is partially occupied or even unoccupied, and metal behaviors indicate a total change of the electronic structure after NO2 adsorption. Besides, the charge transfer between the Sc2CF2 substrate and paramagnetic NO2 also reduces the magnetic moment of the whole system by 0.91 μB.

The rest electronic structure information of the molecules is summarized in Table 1. The bandgap after adsorption does not change much according to Table 1; however, the effective mass changes because of the wave-function mixing. For instance, the orbital of SO2 couples more strongly with the conduction band of Sc2CF2 (figure 3c), and the effective mass of the hole changes by 0.39 m0. Besides, the effective mass of both electrons and holes decreases for each adsorption case, which means the increasing of carrier mobility. It is worth noting that the work function after NO2 adsorption increases by 0.46 eV, which is even larger than what has been reported for graphene.

Furthermore, we explored the optical properties for each system, which may be useful in optical chemistry gas sensors. The imaginary part of dielectric function and adsorption coefficient are plotted at figure 4. At the range of visible light, little change can be found except for a slight intensity decrease and negligible red-shift of the main peaks. However, a clear difference occurs at the ultraviolet range. There are two major absorption peaks (located at 5.3 and 9.4 eV) at the ultraviolet range for pristine Sc2CF2. The intensity of the peak located around 5.3 eV is greatly enhanced after molecules adsorbed at each adsorption system, especially for the case of NO2. Meanwhile, an evident blue-shift happens, and the corresponding peak at the NO2 system moves into around 5.7 eV. The other peak around 9.4 eV changes little except for a slight red-shift. More importantly, the adsorption of each molecule induces a new peak around 9.9 eV, and the NO2 system shows the strongest absorption intensity here. It should also be noted that the adsorption of NO2 induces two additional peaks, which are located around 0.1 and 7.0 eV, respectively. The peak of 0.1 eV has a direct relationship with the partial occupied or unoccupied states near VBM. The quite different optical properties after NO2 adsorption indicate the potential application of Sc2CF2 as a NO2 optical chemistry sensor.

In the aspect of electronic and optical structures, we have found the selectivity of Sc2CF2 to detect NO2. To explore the adsorption thermal stability of NO2 at the Sc2CF2 surface, we performed ab initio molecular dynamics (MD) simulations with an NVT ensemble at a temperature of 300 K for about 11 ps. The time evolution of vertical distances between each atom of NO2 and the top F atom layer of Sc2CF2 is summarized in

![Figure 3](image_url) Band structures and projected DOS of (a) NH3, (b) NO2, (c) SO2, and (d) NO adsorption systems. In the left panels, the gray lines are the band structure of the whole system, and the Fermi level is set as 0. The blue spheres correspond to the states with valid contribution from molecules, and the radii of spheres are proportional to the weight. In the middle and right panels, the dark gray lines are the total DOS of systems, and the blue lines are the PDOS of the molecules.

![Figure 4](image_url) Imaginary part of dielectric function for the polarization vector perpendicular to the surface and absorption coefficient of Sc2CF2 systems with and without molecules adsorbed on the surface.
From the MD simulation, it can be seen that the NO2 molecule becomes 3.18, 3.11, and 2.33 Å at time 5.7 ps, respectively. During the MD process, the vertical distance of N, O-1, and O-2 reaches up to 4.35, 5.49, and 3.72 Å, respectively. However, the NO2 molecule bounces back rapidly, and the vertical distance of N, O-1, and O-2 becomes not so well-saturated, and the adsorption of molecules at gas molecule adsorption systems, and the polar molecules usually show a larger adsorption energy than the nonpolar molecules as we discussed above. Besides, the atoms in the NO2 molecule, and the smallest vertical distance is 3.18 Å, 3.11 Å, and 2.33 Å at time 5.7 ps, respectively. From the MD simulation, it can be seen that the NO2 molecule only vibrates around its equilibrium positions, revealing that NO2 can adsorb on the 2D surface stably at room temperature.

Strain modulation is reported as an effective method to improve the performance of 2D materials. A uniaxial strain can be applied through bending, rolling up, and elongation in experiments, and a homogeneous biaxial strain can also been applied into 2D materials through a substrate that has a different thermal expansion coefficient with 2D materials or a substrate that is a piezoelectric material. Most 2D materials have a fracture strain value over or about 10%, which is a common substrate for 2D material applications in electronics and optoelectronics, which is cheap, flat, and chemically inert. Sc2CF2 layer prefers to contact with O atoms (O-1 or O-2) of the NO2 molecule, and the smallest vertical distance is maintained at the range of 2–4 Å. Compared with the initial structure, the system changes most at time 5.1 ps, when the vertical distance of N, O-1, and O-2 reaches up to 4.35, 5.49, and 3.72 Å, respectively. However, the NO2 molecule bounces back rapidly, and the vertical distance of N, O-1, and O-2 becomes 3.18, 3.11, and 2.33 Å at time 5.7 ps, respectively. From the MD simulation, it can be seen that the NO2 molecule only vibrates around its equilibrium positions, revealing that NO2 can adsorb on the 2D surface stably at room temperature.

Figure 5. (a) Evolution of vertical distances for each atom of NO2 with the top F layer of Sc2CF2 at first-principles molecule dynamics simulations. Snapshot of the NO2 adsorption system at time of (b) 0 ps, (c) 5.1 ps, (d) 6.9 ps, and (e) 11.1 ps, respectively.

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substrate or use h-BN as a substrate on top of SiO2/Si, which will reduce the effect of SiO2.67,68 Here, we considered h-BN as a substrate for Sc2CF2 and explored the effect of h-BN. The lattice constant of h-BN after optimization is 2.41 Å. We used a 3 × 3 Sc2CF2 supercell to match with 4 × 4 h-BN vertically as shown in Figure 7a, and the mismatch is about 1.43%. After optimization, Sc2CF2 is still flat and there exists some buckling in h-BN, which however still maintains the honeycomb structure from the top view. It can be seen from the projected band structure of Sc2CF2/h-BN heterostructures in Figure 7b that the electronic structure of Sc2CF2 can be hardly affected by the h-BN substrate, and a type III band alignment is found. Then we considered four polar polyatomic molecules, that is, NO2, NH3, SO2, and H2O, adsorbed at the surface of Sc2CF2/h-BN. The calculated E_ads and charge transfer from Bader analysis are also plotted in Figure 7c. Here, we used a 3 × 3 supercell, which is a little different from the model before and may have an effect on the value of E_ads and charge transfer. However, the trends of both the adsorption energy and charge transfer are not modulated. The NH3 shows the largest adsorption energy, and the E_ads of NO2, NH3, and SO2 is about 0.21 eV. The largest charge transfer occurs in the case of NO2 and NH3 shows more charge transfer than both H2O and SO2. Furthermore, we calculated the projected band structures for four adsorption systems with the h-BN substrate as shown in Figure S3 in the Supporting Information. In total, the relative position of the molecular orbital with bands of Sc2CF2 changes little, and the LUMO of NO2 is still partially occupied after adsorption, which induces a large amount of charge transfer. These results demonstrate that the h-BN substrate can hardly affect the main results from calculations with a free-standing Sc2CF2. Also note that we considered a pristine Sc2CF2 layer here and the charge transfer will always enhance the carrier density. However, the defect is inevitable in experiments. If the presence of defects makes Sc2CF2 show n-type/p-type conductivity, the carrier density and conductivity do not depend on the charge transfer only. The electron/hole transfer into Sc2CF2 will improve/decrease the n-type conductivity, and the hole/electron transfer into Sc2CF2 will improve/decrease the p-type conductivity. Thus, the predicted obvious selectivity and sensitivity of Sc2CF2 to detect NO2 still need to be confirmed by future experiments.

3. CONCLUSIONS

In summary, we have studied several molecules’ adsorption behaviors on Sc2CF2 and explored the charge transfer between molecules and Sc2CF2. Using first-principles calculations, we have investigated the structural, electronic, and optical properties of adsorption systems. The results show that the largest charge transfer quantity happens between NO2 and Sc2CF2, and optical properties of the system are modified by the NO2 molecule obviously, both of which are due to the special electronic structure, where the LUMO of NO2 is below the VBM of Sc2CF2. Thus, Sc2CF2 can detect the NO2 molecule selectively. Moreover, the MD simulation shows that NO2 can adsorb on the Sc2CF2 surface stably at room temperature. Besides, we explored the effect of biaxial strain on the E_ads and charge transfer quantity of each system, and the result shows that the biaxial strain can enhance both adsorption energy and charge transfer and improve the sensitivity of Sc2CF2 to detect the NO2 molecule. Furthermore, we investigated how h-BN will affect the adsorption behavior and charge transfer, revealing that the h-BN is a potential substrate for Sc2CF2 sensors. Our result predicts that Sc2CF2 can be a promising selective and sensitive sensor to detect the NO2 molecule.

4. COMPUTATIONAL DETAILS

All first-principles calculations were carried out using the Vienna Ab initio Simulation Package (VASP)54 based on DFT with a plane-wave basis set. To treat the exchange–correlation interaction of electrons, we chose the Perdew–Burke–Emzerhof (PBE) functional55 within the generalized gradient approximation (GGA). The electron–ion interactions were described by the projector augmented wave (PAW) potentials.56 The energy cutoff for the plane-wave basis was chosen as 600 eV, and the Brillouin zones (BZs) were sampled by k grids with a uniform spacing of 2π × 0.02 Å−1. Convergence criteria of 10−8 eV for total energy and 0.01 eV/Å for force were adopted for self-consistent calculation and geometry optimization, respectively. We chose a 4 × 4 supercell to simulate the periodic structure of Sc2CF2 monolayers, and a vacuum spacing of 20 Å was added along the direction perpendicular to the 2D sheet (z direction) to avoid interaction between adjacent layers. To properly describe the van der Waals interaction between gas molecules and the Sc2CF2 sheet, we adopted the semi-empirical dispersion-corrected DFT-D3 scheme proposed by Grimme.57 Spin
polarization was included for the systems adsorbed by paramagnetic molecules like NO, NO₂, and O₂. Dipole correction was added in the average potential calculation for the system adsorbed by polar molecules. Various initial configurations with different adsorption sites and molecule orientations of the gas molecules were considered here. The initial distance between the gas molecule and 2D nanosheet was chosen to be 2.0 Å.

**ASSOCIATED CONTENT**

**Supporting Information**

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsomega.1c06027.

DOS of free (a) NO and (b) NO₂ molecule, and inset is the charge density plot (Figure S1); band structures of NO₂ adsorption systems with biaxial strain (Figure S2); and projected band structure with (a) NO₂, (b) NH₃, (c) H₂O, and (d) SO₂ adsorbed at the Sc₂CF₂ surface with h-BN as a substrate. The blue parts are the contribution from the molecule, the gray parts are from Sc₂CF₂, and the orange parts are from h-BN (PDF).

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**Notes**

The authors declare no competing financial interest.

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