Super-PINGU for measuring CP violation

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Abstract

We propose to measure leptonic CP phase, after neutrino mass hierarchy is established, with an upgrade of the PINGU detector and using atmospheric neutrino flux. The upgrade, called super-PINGU, will require a few megaton effective volume at 0.5-1 GeV range to distinguish $\delta$ in the range $\pi/2$-$3\pi/2$ from 0 after 4 years of operation. The distinguishability (similar to significance) of measuring $\delta$ depends crucially on various flux, cross-section, event reconstruction (energy and angle) and flavor identification uncertainties. We explore effects of these uncertainties on the distinguishability of measuring CP phase and suggest possible ways to minimize their impact.

Keywords: Neutrino oscillation, CP violation, atmospheric neutrino

1. Introduction

Establishing CP violation in the leptonic sector is an outstanding problem in particle physics. Atmospheric neutrino flux measurements in large water/ice detectors can be used to determine the Dirac CP phase $\delta$. Information on different $\delta$ values is encoded in the neutrino oscillation probabilities after propagation inside the earth (matter effect). A systematic shift of the probabilities with increasing $\delta$ in a wide, $\sim 0.2$--2 GeV, energy range is key to measure $\delta$ with atmospheric neutrino flux \cite{1}.

It has been found recently that the PINGU and ORCA detector with $\sim 3$ GeV threshold will have good sensitivity to determining the neutrino mass hierarchy \cite{2,3,4,5}. However, measurement of $\delta$ would require a detector with larger effective volume and improved characteristics in the < 2 GeV range and in this context a future upgrade of PINGU (and also of ORCA), called Super-PINGU, was proposed in Ref. \cite{1} with detailed estimate of sensitivity. Here we present highlights from that work.

2. Methodology and Results

To measure CP phase with atmospheric neutrinos, we calculate neutrino events in Super-PINGU by varying $\delta$ and compare with $\delta = 0$, keeping all other oscillation parameters fixed. We use an effective mass (both for $\nu_\mu$ and $\nu_e$) parameterized as

$$\rho V_{\text{eff}}(E_\nu) = 2 \pi N_A \rho T \int_{\Delta E_i} \int_{\Delta \cos \theta_z} dE_\nu V_{\alpha \beta}(E_\nu) d\alpha(E_\nu, \theta_z).$$

which can be realized for a total of 126 strings and 60 DOMs per string. This gives an effective mass of $\sim 2.8$ Mt at $\sim 1$--2 GeV, which is 4 times larger than PINGU. The number of neutrino events for a particular flavor $\alpha = e, \mu$ with energies and zenith angles in small bins $\Delta(E_i)$ and $\Delta(\cos \theta_z)$ marked by subscript $ij$ can be calculated as

$$N_{ij,\alpha} = 2 \pi N_A \rho T \int_{\Delta E_i} \int_{\Delta \cos \theta_z} dE_\nu V_{\alpha \beta}(E_\nu) d\alpha(E_\nu, \theta_z).$$

Here $T$ is the exposure time, $N_A$ is the Avogadro’s number. The density of events of type $\alpha$, $d_\alpha$ (the number of events per unit time per target nucleon), is given by

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rection using PINGU reconstruction functions with a factor \( \nu_0 \) for neutrino energy and direction after smearing of the true energy and direction. The normal mass hierarchy is assumed.

Figure 1: Distributions of \( \delta \) events in the \( \nu_\mu + \bar{\nu}_\mu \) plane. \( E_\nu \) and \( \theta_2 \) are reconstructed neutrino energy and direction after smearing of the true energy and direction using PINGU reconstruction functions with a factor \( \sqrt{3} \) better resolutions. Normal mass hierarchy is assumed.

\[
d_\alpha(E_\nu, \theta_2) = \sigma_\alpha \Phi_\alpha + \sigma_\beta \Phi_\beta \quad \text{in terms of fluxes at the detector, and} \quad \Phi_\alpha = \Phi^0_\mu P^0_\mu + \Phi^0_\nu P^0_\nu \quad \text{with corresponding oscillation probabilities} \quad P_{\mu\alpha} \quad \text{and} \quad P_{\nu\alpha} \ . \]

The original muon and electron neutrino fluxes at the production are \( \Phi^0_\mu = \Phi^0_\mu(E_\nu, \theta_2) \) and \( \Phi^0_\nu = \Phi^0_\nu(E_\nu, \theta_2) \).

We have computed the distributions of \( \nu_\mu + \bar{\nu}_\mu \) and \( \nu_e + \bar{\nu}_e \) events for \( \delta = 0 \) and \( \delta \neq 0 \) and take difference between the distributions in the \( E_\nu - \cos \theta_2 \) plane to study their properties. Since there are errors associated with reconstructing the true neutrino energy and directions, we smear the ideal distributions with the energy and angular resolution functions of the detector to mimic the real situation. To estimate the sensitivity of measuring a CP phase different from zero we employ a distinguishability parameter defined as

\[
S_{ij}(f) = (N^0_{ij} - N^f_{ij})/\sigma_{ij}(f),
\]

where \( N^0_{ij} \) and \( N^f_{ij} \) are the reconstructed number of events in the \( ij \)-th bin in the \( E_\nu - \cos \theta_2 \) plane for \( \delta \) and \( \delta = 0 \), respectively, and \( \sigma^2_{ij}(f) = N^0_{ij} + (f N^0_{ij})^2 \) is the total error in the \( ij \)-th bin. Parameter \( f \) is a measure of uncorrelated systematic errors \( \xi \) and we take \( f = 2.5\% \).

The total distinguishability

\[
S_\sigma = \sqrt{\sum_{ij} S_{ij}^2}
\]

is a quick measure of significance.

Figure 1 shows the distinguishability \( S_{ij} \) for \( \nu_\mu + \bar{\nu}_\mu \) events with \( f = 0 \) for different \( \delta \) values and using 1-year of Super-PINGU data. Normal mass hierarchy is assumed. The shape of the distributions, specially their domain structures, is largely explained as due to grids of solar, atmospheric and interference magic lines in the \( E_\nu - \cos \theta_2 \) plane. The oscillation probabilities are independent of \( \delta \) along these lines, thus separating regions of same sign distinguishability. Figure 2 shows \( S_{ij} \) distributions for \( \nu_e + \bar{\nu}_e \) events.

The uncertainties associated with atmospheric neutrino flux, \( \nu N \) cross section, effective volume, etc. affect neutrino event distributions across bins in the \( E_\nu - \cos \theta_2 \) plane. We include effects of these correlated uncertainties in our calculation with analogy to the pull method in \( \chi^2 \) analysis. In particular we minimize the following distinguishability parameter

\[
S_{\sigma l}^{\xi} = \left[ \sum_{i \in \xi} \sum_{ij} \frac{[N^0_{ij}(\xi_l) - N^0_{ij}(\xi_l^\ast)]^2}{\sigma^2_{ij}} \right]
\]
\[
\sum_{k} \left( \frac{(\xi_k - \xi_{st}^k)^2}{\sigma_k^2} \right)^{1/2},
\]

where \(\xi_k\) are the pull variables and \(\xi_{st}^k\) are their standard values. The event distributions with varying \(\xi_k\) are calculated as

\[
N_{ij}^\delta(\xi_k) = \alpha z_l(E/2 \text{ GeV})^\eta \times [1 + \beta(0.5 + \cos \theta_z)] N_{ij}^\delta(\xi_{st}^k),
\]

where \(\alpha\) is the overall normalization factor with the error \(\sigma_\alpha = 0.2\), \(z_l\) is the flux (flavor) ratio uncertainty (\(z_e \equiv 1\) for \(\nu_e\) events), with the error \(\sigma_{z_l} = 0.05\); \(\eta\) is the energy tilt parameter with \(\sigma_\eta = 0.1\); \(\beta\) is the zenith angle tilt with \(\sigma_{\beta} = 0.04\). Figure 3 shows the \(S_{tot}^\delta\) minimized over \(\xi_k\) for different correlated uncertainties as well as for no correlated uncertainties. A threshold energy of 0.5 GeV has been assumed. Note that the contributions of \(\nu_e\) and \(\nu_\mu\) channels to \(S_{tot}^\delta\) are comparable.

3. Discussion

We estimate that after 4 years of operation and 2.5% systematics, Super-PINGU with 0.5 GeV threshold will be able to distinguish \(\delta = \pi/4, \pi/2, \pi, 3\pi/2\) from zero with \(S_{tot}(\pi/4) = (1 - 3), S_{tot}(\pi/2) = (3 - 8), S_{tot}(\pi) = (6 - 14), S_{tot}(3\pi/2) = (3 - 8)\). The ranges depend on effects of different correlated systematics. These values are a factor 4–6 improvement over the sensitivity of PINGU to \(\delta\) with 3 GeV threshold.

The sensitivity of Super-PINGU to \(\delta\) can be further improved with following possibilities:

- Decrease of energy threshold to 0.2 GeV from 0.5 GeV with a denser array. This may increase sensitivity by 30%.
- Stringent kinematical cut can be used to create a high-quality event sample with better reconstruction of the neutrino energy, direction and flavor.
- An increased exposure time will also increase the sensitivity to CP by a factor \(\propto \sqrt{t}\).
- Improved flavor identification at low energies.
- Increase the density of DOMs or photocathode coverage.
- Statistical separation between the neutrino and antineutrino events.

Our results show that Super-PINGU can be competitive to other proposals for measuring the leptonic CP phase associated with long base-line (LBL) accelerator experiments.

References

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Figure 3: Distinguishability of measuring CP phase in the \(\nu_\mu\) and \(\nu_e\) channels, with effects of various correlated uncertainties. 2.5% uncorrelated uncertainty has been assumed in all cases.