Realistic Hi scale heights of Milky Way-mass galaxies in the FIREbox cosmological volume

Jindra Gensior, 1* Robert Feldmann, 1 Lucio Mayer, 1 Andrew Wetzel, 2 Philip F. Hopkins, 3 and Claude-André Faucher-Giguère 4

1 Institute for Computational Science, University of Zürich, Winterthurerstrasse 190, 8057 Zürich, Switzerland
2 Department of Physics & Astronomy, University of California, Davis, CA 95616, USA
3 California Institute of Technology, TAPIR, Mailcode 350-17, Pasadena, CA 91125, USA
4 CIERA and Department of Physics and Astronomy, Northwestern University, 1800 Sherman Ave, Evanston, IL 60201, USA

Accepted XXX. Received YYY; in original form ZZZ.

ABSTRACT

Accurately reproducing the thin cold gas discs observed in nearby spiral galaxies has been a long standing issue in cosmological simulations. Here, we present measurements of the radially resolved Hi scale height in 22 non-interacting Milky Way-mass galaxies from the FIREbox cosmological volume. We measure the Hi scale heights using five different approaches commonly used in the literature: fitting the vertical volume density distribution with a Gaussian, the distance between maximum and half-maximum of the vertical volume density distribution, a semi-empirical description using the velocity dispersion and the galactic gravitational potential, the analytic assumption of hydrostatic equilibrium, and the distance from the midplane which encloses ≥60 per cent of the Hi mass. We find median Hi scale heights, measured using the vertical volume distribution, that range from ~100 pc in the galactic centres to ~800 pc in the outskirts and are in excellent agreement with recent observational results. We speculate that the presence of a realistic multiphase interstellar medium, including cold gas, and realistic stellar feedback are the drivers behind the realistic Hi scale heights.

Key words: galaxies: ISM – galaxies: spiral – galaxies: structure – ISM: structure – ISM: kinematics and dynamics

1 INTRODUCTION

Producing realistic disc galaxies in cosmological scale simulations has been a long-standing problem for numerical astrophysics (e.g. Naab & Ostriker 2017, and the references therein). Only in the past decade have cosmological zoom-in simulations and volumes been able to reproduce the rotation curves and thin stellar discs seen in late-type spiral galaxies (e.g. Guedes et al. 2011; Stinson et al. 2013; Aumer et al. 2013; Hopkins et al. 2014; Agerzt & Kravtsov 2015; Grand et al. 2017; Ma et al. 2017; Garrison-Kimmel et al. 2018; Pillepich et al. 2019; Buck et al. 2020). This success if often attributed to the (stellar) feedback models used, as the thickness of the stellar disc is sensitive to the feedback strength (e.g. Sokolowska et al. 2017).

By contrast, recent cosmological (zoom) simulations which investigated the H I content of L* galaxies found the H I discs to be more thick. H I scale heights of the Auriga (Grand et al. 2017) galaxies within the optical radius range from ~1–10 kpc (Marinacci et al. 2017). Similarly, the EAGLE (Schaye et al. 2015) galaxies have H I discs with scale heights > 1.5 kpc, which Bahé et al. (2016) interpreted as consequence of too strong feedback.

The scale height or thickness of gas discs is relevant for determining their susceptibility to gravitational instability (e.g. Romeo & Wiebert 2011). Whether a (super)bubble driven by supernovae (SNe) can break out of the galactic disc and pollute the circumgalactic medium (CGM) with metals, or whether it stalls and deposits all ejecta in the gas disc of the galaxy also depends on the scale height of the cold gas (e.g. Silich & Tenorio-Tagle 2001; Orr et al. 2021, 2022). Thus the scale height has direct implications for the star formation and feedback cycle of galaxies, as well as the metal pollution of the CGM. In turn, feedback will affect the H I structure (e.g. Bahé et al. 2016). Therefore, the H I scale height can be a sensitive test for sub-grid physics, as it is unclear a priori whether well validated sub-grid models like Fire-2 will be able to reproduce observed scale heights. Accurately modelling (spiral) galaxies and their H I gas content is also highly relevant with respect to the upcoming Square Kilometer Array (SKA; Dowden et al. 2009) observatory and its precursors. Both to make higher fidelity predictions, and to be able to better interpret observational results.

Encouragingly, thin cold gas discs have been found in sev-
eral $L_*$ galaxies simulated with the FIRE-2 model (Sanderson et al. 2020; Gurvich et al. 2020; Trapp et al. 2022), but no systematic exploration with a focus on the HI has been conducted yet. The FIREbox cosmological volume (Feldmann et al. 2022), with its high spatial resolution ($\sim 20$ pc), and multiphase ISM, that yields 20–30 $L_*$ galaxies, is ideally suited to systematically analyse the HI discs of $L_*$ galaxies. In this letter, we focus on their scale heights in particular, measuring them with the variety of different methods used in the literature.

2 FIREBOX
In this letter we study galaxies drawn from the $(22.1$ Mpc)$^3$ FIREbox cosmological volume simulation (Feldmann et al. 2022) that is part of the Feedback In Realistic Environments (FIRE) project$^1$. The simulation was run with the meshless-finite-mass code GIZMO$^2$ (Hopkins 2015) and the FIRE-2 (Hopkins et al. 2018) sub-grid physics. Specifically, these simulations use the Hopkins et al. (2014) cooling and heating rates valid for temperatures ranging from $10^{10}$–$10^{12}$ K, including heating and photoionisation from a Faucher-Giguère et al. (2009) UV Background, naturally resulting in a multiphase interstellar medium (ISM) that includes a cold gas phase. Stars form with a local efficiency of 100 per cent from gas that is molecular, self-gravitating, Jeans unstable and above a density threshold of $n \geq 300$ cm$^{-3}$. A variety of stellar feedback channels are included, namely SN Type II and Ia, stellar winds from massive OB and evolved AGB stars, as well as photoionization, photoelectric heating and radiation pressure.

FIREbox has a mass resolution of $m_\text{b} = 6.3 \times 10^4$ M$_\odot$ ($m_{\text{DM}} = 3.3 \times 10^5$ M$_\odot$), fixed gravitational softening for the stars (12 pc) and dark matter (80 pc), and adaptive gravitational softening for the gas down to a softening length of 1.5 pc (all in proper units). Galaxies were identified at $z = 0$ using the AMIGA halo finder (Gill et al. 2004; Knollmann & Knebe 2009). We restrict our analysis to centrals in the virial mass range $11.85 < \log_{10}(M_{\text{vir}}/M_\odot) < 12.48$, such that the mean of the resultant 27 galaxies matches the Milky Way $M_{\text{vir}} = 1.3 \pm 0.1 \times 10^{12}$ from Bland-Hawthorn & Gerhard (2016). To better match the observational samples, we visually inspect the galaxies and exclude those that are interacting or undergoing a merger. The stellar masses of our final sample of 22 galaxies range from $3.5 \times 10^{10}$ to $1.6 \times 10^{11}$ M$_\odot$, in good agreement with the mass range of the BLUEDISK galaxies and the higher mass THINGS galaxies studied by Patra (2020), while Bacchini et al. (2019) also include some lower mass THINGS galaxies.

To obtain the HI mass fraction of each gas particle, we first calculate the H$_2$ fraction of the neutral gas following Krumholz & Gnedin (2011), then subtract the molecular fraction, as well as contributions from metals and Helium, from the neutral gas fraction of each particle.

3 MEASURING THE HI SCALE HEIGHT
The scale height of real galaxies can only be measured directly for edge-on galaxies, where the vertical gas distribution can be fit with a Gaussian (e.g. Yim et al. 2011). Below, we briefly discuss the different ways the HI scale height can be measured for galaxies at lower inclinations.

Assuming that the gas disc is in hydrostatic equilibrium, with turbulent pressure$^3$ balancing the effect of gravity, it is possible to derive an analytic estimate of the HI scale height (see Appendix A of Bacchini et al. 2019, and references therein).

$$h_{\text{HI}}(R) \approx \sigma_{\text{HI}}(R)/\sqrt{4\pi G \left(\rho(R,0) - \frac{1}{2\pi G} \frac{v_c(R)}{R} \frac{\partial v_c(R)}{\partial R}\right)},$$

where $h_{\text{HI}}$ is the scale height of the HI disc, $\sigma_{\text{HI}}$ is the mass-weighted line-of-sight velocity dispersion in the direction perpendicular to the disc plane, $v_c$ is the circular velocity and $\rho(R,0)$ is the mean mass-weighted density at the galaxy mid-plane, all as a function of galactocentric radius $R$. We calculate all quantities as running mean in overlapping radial bins of 1 kpc width.

The analytic approximation neglects the contribution of the gas self-gravity to the gravitational potential. To include gas self-gravity, and thus determine $h_{\text{HI}}$ more accurately, one needs to solve for the volume density numerically. It is done by solving the equation of the hydrostatic equilibrium condition under the constraints of the rotation curve and velocity dispersion, iterating until the so derived estimates converge. This method has been used in two subtly-different ways in the literature.

The first approach assumes that the vertical volume density profile has a Gaussian shape that depends on the HI scale height, i.e. $\rho(z) \propto \exp(-z^2/(2h_{\text{HI}}^2))$. It was used by e.g. Bacchini et al. (2019) to measure the HI scale heights of 12 THINGS galaxies. To mimic this method, we fit the HI mass-weighted vertical volume density profile of gas with $z \leq 2$ kpc in radial annuli of 1 kpc width, with a Gaussian function and radiation pressure.

With the half-width half maximum (HWHM) approach, no functional form is assumed for the HI volume density. The HI scale height is defined as the distance between the midplane (maximum volume density) and one of the points where the HI volume density reaches its half-maximum. It was used by e.g. Patra (2020) to measure $h_{\text{HI}}$ in 9 THINGS galaxies. We determine this distance for the HI mass-weighted vertical volume density profile of gas with $z \leq 2$ kpc in radial annuli of 1 kpc width.

Another way to estimate $h_{\text{HI}}$ using a semi-empirical formula was originally derived by Wilson et al. (2019), assuming hydrostatic equilibrium but taking into account observational results regarding the gravitational potentials of galaxies. We follow the parametrisation of Randriamampandry et al. (2021), who used this equation to measure $h_{\text{HI}}$ for the BLUEDISK galaxies. The HI scale height can be expressed as:

$$h(R) = \left[\sigma_{\text{HI},0}^2(R) / \pi G \Sigma_{\text{total}} \right]^{1/2} \left[1 + \sigma_{\text{HI},0}^2(R) / 2v_{\text{max}}^2 \left(\frac{M_{\text{dyn}}}{M_{\text{HI}} + M_*}\right)^2\right]^{-1},$$

where $\sigma_{\text{HI},0}$ is the central HI velocity dispersion, $\Sigma_{\text{total}}$ is the total surface density within the gas layer, $v_{\text{max}}$ is the peak of the galaxy’s rotation curve, $M_{\text{dyn}}$ is the total (dynamical) mass, $M_{\text{HI}}$ the total mass of HI and $M_*$ the total stellar mass.

Finally, Marinacci et al. (2017) measure the HI scale height as the distance from the midplane where the HI mass enclosed within a cylinder spanning $\pm 30$ kpc from the midplane of the galaxy equals $3$ Bacchini et al. (2019) explicitly assume the gas to be vertically isothermal. However, it has also been shown in numerical simulations with effective feedback that the turbulent pressure dominates over the thermal pressure (Benincasa et al. 2016).
4 PROPERTIES OF FIREBOX GALAXIES

We show the radial profiles of several galaxy properties (HI velocity dispersion, HI surface density, HI mass-weighted volume density at the midplane and rotation velocity curves), relevant for calculating the HI scale height either analytically (Equation 1) or using the Wilson et al. (2019) semi-empirical equation (2), in Figure 1. The median \( \sigma_{HI} \) radial profile peaks at the centre of the galaxy with 40 km s\(^{-1}\), declining with increasing galactocentric radius to a value of \( \sim 17 \) km s\(^{-1}\) for \( R < 15 \) kpc, and \( \sim 27 \) km s\(^{-1}\) for \( R > 15 \) kpc. The median consistently lies above the often assumed canonical value of \( \sigma_{HI} \sim 10 \) km s\(^{-1}\) (e.g. Tamburro et al. 2009), also seen in the BLUEDISK galaxies (Randriamampandry et al. 2021). However, the HI volume density for the Firebox galaxies is calculated for all HI gas, including the warm component, for which Janamasimanana et al. (2012) found an average velocity dispersion of 16.8 km s\(^{-1}\) in good agreement with the median \( \sigma_{HI} \) at intermediate radii. The rise of the median velocity dispersion in the central 5 kpc is indicative of shear-driven turbulence induced by a bulge (see e.g. Figure 9 in Gensior et al. 2020). The median rotation curve in the top right panel shows the characteristic bump expected from a spheroidal component around \( R = 2 \) kpc, supporting this interpretation. The velocity dispersion profiles of the THINGS galaxies with bulges studied in Bacchini et al. (2019) and Patra (2020) qualitatively match the radial profiles in Figure 1, although the central velocity dispersions of the THINGS galaxies reach \( \sim 30 \) km s\(^{-1}\) at most. The HI surface density is calculated as the HI mass enclosed in overlapping radial annuli of width 1 kpc, divided by the area of the annulus. The median \( \Sigma_{HI} \) stays approximately constant at \( \sim 7.5 \) M\(_{\odot}\)pc\(^{-2}\) out to a galactocentric radius of 8 kpc, before declining to a density \( \sim 1 \) M\(_{\odot}\)pc\(^{-2}\) at a median HI disc radius \( R_{HI} \approx 20\) kpc, in good agreement with observed \( \Sigma_{HI} \) profiles (e.g. Leroy et al. 2008; Wang et al. 2014; Randriamampandry et al. 2021). The median HI mass-weighted volume density at the midplane peaks at \( \sim 18 \) cm\(^{-3}\) in the galactic centre and declines steeply with increasing galactocentric radius, falling below \( 1 \) cm\(^{-3}\) at a radius of 13 kpc.

5 HI SCALE HEIGHT OF FIREBOX GALAXIES

Figure 2 shows the HI scale height of the 22 \( L_\star \) Firebox galaxies measured using the different approaches outlined in Section 3. The top three and bottom two leftmost panels show the measurements for individual galaxies, with median, error on the median determined via bootstrapping, and 16th-to-84th percentile variation for the sample indicated by the coloured line and dark and light shaded regions respectively, compared with the observational result using the same method where applicable. The bottom right panel shows a compilation of the median measurements from the different methods, simulations and observations.

All median HI scale height profile obtained from the simulated galaxies show disc flaring, i.e. follow the trend of increasing \( h_{HI} \) with increasing galactocentric radius. However, the different methods vary by more than an order of magnitude in their estimate of the central \( h_{HI} \), the amount of disc flaring predicted, as well as the radial gradient of the profile.

Using the vertical profile of the HI mass-weighted volume density to determine the \( h_{HI} \), either by fitting the gas distribution with a Gaussian, or calculating the HWHM, yields results in excellent agreement with measurements from the THINGS galaxies that use the same methodology (Bacchini et al. 2019; Patra 2020). The HWHM method finds a lower central scale height compared to the Gaussian fit (60 and 85 pc respectively), and predicts a steeper increase to a scale height of 500 pc at \( R = 8 \) kpc, before oscillating and gradually increasing to 900 pc at \( R = 25 \) kpc. The HI scale height obtained from the Gaussian fit increases more shallowly with \( R \), reaching 800 pc at \( R = 14 \) kpc, staying nearly constant at larger radii (reaching 1 kpc at \( R = 25 \) kpc). The analytic determination of \( h_{HI} \) agrees well with the volume density estimates and the observations for \( R < 15 \) kpc, but predicts too much flaring at larger galactocentric radii, likely driven by the increase of \( \sigma_{HI} \) beyond this radius.

Determining the HI scale height from Equation 2 yields a median central scale height of \( 13 \) pc, increasing to \( 100 \) pc at a galactocentric radius of 12 kpc, and flaring further until reaching a scale height of \( \approx 0.9 \) kpc at \( R > 17.5 \) kpc. The median of the simulated galaxies lies significantly below the median Randriamampandry et al. (2021) BLUEDISK measurement, that only shows a very modest increase from 200 to 350 pc, for \( R < 12 \) kpc. Due to the large flaring predicted, the median of the simulated galaxies lies within the 16th-to-84th percentile of BLUEDISK at larger radii. The large discrepancies at small galactocentric radii are a consequence of the strong dependence of \( h_{HI} \) on \( \sigma_{HI} \) in Equation 2, as the central velocity dispersions of the Firebox \( L_\star \) galaxies far exceed the \( \sim 10 \) km s\(^{-1}\) of the BLUEDISK galaxies\(^4\). At large galactocentric radii the strong flaring is driven by the increase in the median \( \sigma_{HI} \) from 17 to 27 km s\(^{-1}\) (see the top left panel of Figure 1). Defining the scale height as the distance from the midplane that encloses 63 per cent of the total HI mass in a hollow cylinder of height 60 kpc systematically overpredicts \( h_{HI} \) compared to all other methods of obtaining a scale height, with a median scale height of \( \sim 18 \) cm\(^{-3}\) in the galactic centre and declines steeply with increasing galactocentric radius, falling below 1 cm\(^{-3}\) at a radius of 13 kpc.

\(^4\) However, the BLUEDISK observations have coarse resolution and lack \( \sigma \) measurements in their centres, making it difficult to determine how much of the discrepancy between the BLUEDISK observations and simulations is due to (potential) morphological differences or strong FIRE-2 feedback potentially further enhancing the velocity dispersion.
Figure 2. Radial profiles of the H\textsubscript{i} scale heights of the Firebox $L_*$ galaxies estimated from Equation 1, 2, a Gaussian fit to the vertical volume density profile, distance from the galaxy mid-plane that encloses 63 per cent of the H\textsubscript{i} mass and the vertical distance between the maximum and half of the maximum volume density, from top left to bottom left respectively, with observations from Randriamampandry et al. (2021), Bacchini et al. (2019), and Patra (2020) over-plotted, where the same method of estimating the H\textsubscript{i} scale height is used. Coloured lines indicate the median, the dark shaded regions the error on the median determined via bootstrapping and the light shaded regions the 16th-to-84th percentile variation of the data. The bottom right panel combines the median measurements of observations and simulations presented in the other five panels. The H\textsubscript{i} scale heights measured from the vertical volume density distribution are in excellent agreement with observational measurements.

350 pc at the galactic centre, increasing to 10 kpc by a galactocentric radius of 25 kpc. The total H\textsubscript{i} mass enclosed depends on the height of the cylinder, thus using a smaller distance from the midplane would yield lower scale height estimates.

6 WHY DOES FIREBOX PRODUCE REALISTIC H\textsubscript{i} SCALE HEIGHTS?

We can speculate what makes Firebox so successful at producing the thin H\textsubscript{i} discs shown in Figure 2 based on possible explanations put forward in the literature. In some cases (e.g. Guedes et al. 2011; Pillepich et al. 2019) high resolution has been crucial to reproducing thin (stellar) discs. However, as we show in Appendix A, using the vertical volume density profile to determine the H\textsubscript{i} scale height yields good agreement for simulations differing by nearly two orders of magnitude in spatial resolution (minimum gas softening ranging from 1.5 to 100 pc). This suggests that the other modelling choices are more relevant for producing realistic galaxy discs. However, the minimum gravitational softening likely needs to be lower than the scale height, to reproduce thin scale heights. This is consistent with the Firebox results, as the lowest resolution re-run has a minimum gravitational force softening below the smallest measured $h_{\text{HI}}$ of $\sim$100 pc.

Trayford et al. (2017) suggest that using an effective equation of state instead of directly modelling the multiphase ISM leads to too thick stellar discs in the EAGLE galaxies. Pillepich et al. (2019) also report a 0.5 dex variation in scale height when adjusting the parameters of their equation of state model. The Fire-2 model includes cooling down to temperatures of 10 K, thus we can self-consistently model a multiphase ISM including cold gas, plausibly a major factor in producing realistic gas scale heights.

Many authors have argued that the inclusion of more sophisticated feedback models such as superbubble feedback (e.g. Keller et al. 2015) or early stellar feedback are key to produce realistic $L_*$ galaxies (e.g. Stinson et al. 2013; Aumer et al. 2013; Hopkins et al. 2014). This suggests that the early stellar feedback implemented in Fire-2 could be crucial to producing thin gas discs, by pre-processing the ISM gas, thus enabling the effective driving of localised SNe outflows even at coarse mass resolution (see Figure A1), as opposed to SNe depositing more of their energy and momentum within the ISM, increasing the local pressure and thereby increasing the overall scale height. Apart from dedicated new simulations, it might be informative to see whether other simulations with a cold ISM but no early stellar feedback (e.g. EMP-Pathfinder; Reina-Campos et al. 2022) and those with early stellar feedback (e.g. NIHAO; Buck et al. 2020) also show thin H\textsubscript{i} gas discs.

7 CONCLUDING REMARKS

We measure the H\textsubscript{i} scale height of 22 Milky Way-mass galaxies from the Firebox pathfinder volume, using five different methods. We find
that the two methods using the vertical volume density distribution (Gaussian fit and HWHM) are in excellent agreement with each other and the observations. We speculate that the combination of detailed modelling of the cold ISM and early stellar feedback are responsible for these realistically thin H\textsc{i} gas discs, a trend that is independent of resolution.

The analytic hydrostatic equilibrium assumption works well for the central ~15 kpc, but overpredicts the scale height compared to observations and measurements from the vertical volume density distribution at large radii, due to the rising velocity dispersion. The Wilson et al. (2019) equation underpredicts the scale height in the central region of the galaxies, due to the high central velocity dispersion of the Firebox galaxies, but is in good agreement with the volume density estimates at large radii. Using the H\textsc{i} mass enclosed as an estimate consistently overpredicts the gas scale height at all radii, although the general trend of disc flaring at large radii is present as well.

Turbulent theory predicts that the H\textsc{i} scale height is a transition scale, where turbulence changes from three-dimensional turbulence on small scales to two-dimensional turbulence on larger scales (e.g. Agertz & Kravtsov 2015). This could cause a break in the H\textsc{i} power spectra of galaxies (e.g. Dutta et al. 2008). Since Firebox produces galaxies with realistic H\textsc{i} scale heights, it is well suited to study H\textsc{i} turbulence and its drivers and other processes that depend on the scale height. Futhermore, this makes Firebox a great resource for more in-depth studies of the H\textsc{i} content of galaxies in general, and for making predictions for upcoming observatories such as the SKA.

**ACKNOWLEDGEMENTS**

JG, RF and LM gratefully acknowledge financial support from the Swiss National Science Foundation (grant no CRSII5_193826). RF acknowledges financial support from the Swiss National Science Foundation (grant no PP00P2_194814 and 200021_188552). AW received support from: NSF via CAREER award AST-2045928 and grant AST-2107772; NASA ATP grant 80NSSC20K0513; HST grants AR-15809, GO-15902, GO-16273 from STScI. CAFG was supported by NSF through grants AR-15809, GO-15902, GO-16273 from STScI. CAFG was supported by NSF via CAREER award AST-2045928 and grants TM2-23005X; and by the Research Corporation for Science Advancement through a Cottrell Scholar Award. We acknowledge PRACE for awarding us access to MareNostrum at the Barcelona Supercomputing Center (BSC), Spain. This work was supported in part by a grant from the Swiss National Supercomputing Centre (CSCS) under project IDs s697 and s698. We acknowledge access to Piz Daint at the Swiss National Supercomputing Centre, Switzerland under the University of Zurich’s share with the project ID uzh18. This work made use of infrastructure services provided by S3IT (www.s3it.uzh.ch), the Service and Support for Science IT team at the University of Zurich.

**DATA AVAILABILITY**

The data underlying this article will be shared on reasonable request to the corresponding author.

**REFERENCES**

Agertz O., Kravtsov A. V., 2015, ApJ, 804, 18

Aumer M., White S. D. M., Naab T., Scannapieco C., 2013, MNRAS, 434, 3142

Bacchini C., Fraternali F., Iorio G., Pezzulli G., 2019, A&A, 622, A64

Bahé Y. M., et al., 2016, MNRAS, 456, 1115

Benincasa S. M., Wadsley J., Couchman H. M. P., Keller B. W., 2016, MNRAS, 462, 3053

Bland-Hawthorn J., Gerhard O., 2016, ARA&A, 54, 529

Buck T., Obreja A., Macciò A. V., Minchev I., Dutton A. A., Ostriker J. P., 2020, MNRAS, 491, 3461

Dewdney P. E., Hall P. J., Schilizzi R. T., Lazio T. J. L. W., 2009, IEEE Proceedings, 97, 1482

Dutta P., Begum A., Bharadwaj S., Chengalur J. N., 2008, MNRAS, 384, L34

Faucher-Giguère C.-A., Lidz A., Zaldarriaga M., Hernquist L., 2009, ApJ, 703, 1416

Feldmann R., et al., 2022, arXiv e-prints, p. arXiv:2205.15325

Garrison-Kimmel S., et al., 2018, MNRAS, 481, 4133

Gensler J., Kuijken J. M. D., Keller B. W., 2020, MNRAS, 495, 199

Gill S. P. D., Knebe A., Gibson B. K., 2004, MNRAS, 351, 399

Grand R. J. J., et al., 2017, MNRAS, 467, 179

Guedes J., Callegari S., Madau P., Mayer L., 2011, ApJ, 742, 76

Gurvich A. B., et al., 2020, MNRAS, 498, 3664

Hopkins P. F., 2015, MNRAS, 450, 53

Hopkins P. F., Kereš D., Ohtobe J., Faucher-Giguère C.-A., Quataert E., Murray N., Bullock J. S., 2014, MNRAS, 445, 581

Hopkins P. F., et al., 2018, MNRAS, 480, 800

Ianjamasimanana R., de Blok W. J. G., Walter F., Heald G. H., 2012, AJ, 144, 96

Keller B. W., Wadsley J., Couchman H. M. P., 2015, MNRAS, 453, 3499

Knollmann S. R., Knebe A., 2009, ApJS, 182, 608

Krumholz M. R., Gnedin N. Y., 2011, ApJ, 732, 729

Leroy A. K., Walter F., Brinks E., Bigiel F., de Blok W. J. G., Madore B., Thornley M. D., 2008, AJ, 136, 2782

Ma X., Hopkins P. F., Wetzel A. R., Kirby E. N., Anglés-Alcázar D., Faucher-Giguère C.-A., Quataert E., Murray N., Bullock J. S., 2014, MNRAS, 447, 2430

Marinacci F., Grand R. J. J., Pakmor R., Springel V., Gómez F. A., Frenk C. S., White S. D. M., 2017, MNRAS, 466, 3859

Naab T., Ostriker J. P., 2017, ARA&A, 55, 59

Orr M. E., Fielding D. B., Hayward C. C., Burkhardt B., 2021, arXiv e-prints, p. arXiv:2109.14656

Patra N. N., 2020, MNRAS, 499, 2063

Pillepich A., et al., 2019, MNRAS, 490, 3196

Randriamampandry T. H., Wang J., Mogotsi K. M., 2021, ApJ, 916, 26

Reina-Campos M., Keller B. W., Kuijken J. M. D., Gensler J., Travijilo-Gomez S., Jefferson S. M. R., Pfeffer J. L., Sills A., 2022, arXiv e-prints, p. arXiv:2202.06961

Romeo A. B., Wiepert J., 2011, MNRAS, 416, 1191

Sandorson R. E., et al., 2020, ApJS, 246, 6

Schaye J., et al., 2015, MNRAS, 446, 521

Silich S., Tenorio-Tagle G., 2001, ApJ, 552, 91

Sokolowska A., Capelo R. P., Fall S. M., Mayer L., Shen S., Bonoli S., 2017, ApJ, 835, 289

Stinson G. S., Brook C., Macciò A. V., Wadsley J., Quinn T. R., Couchman H. M. P., 2013, MNRAS, 428, 129

Tamburro D., Rix H. W., Leroy A. K., Mac Low M. M., Walter F., Kennicutt R. C., Brinks E., de Blok W. J. G., 2009, AJ, 137, 4424

Trapp C. W., et al., 2022, MNRAS, 509, 4149

Trayford J. W., et al., 2017, MNRAS, 470, 771

Walter F., Brinks E., de Blok W. J. G., Bigiel F., Kennicutt Robert C. J., Thornley M. D., Leroy A. 2008, AJ, 136, 2563

Wang J., et al., 2013, MNRAS, 433, 270

Wang J., et al., 2014, MNRAS, 441, 2159

Wilson C. D., Elmegreen B. G., Binami E., Brunetti N., 2019, ApJ, 882, 5

Yim K., Wong T., Howk J. C., van der Halst J. M., 2011, AJ, 141, 48

Yim K., Wong T., Xue R., Rand R. J., Rosolowsky E., van der Halst J. M., Benjamin R., Murphy E. J., 2014, AJ, 148, 127
APPENDIX A: RESOLUTION TEST

We study the effect of spatial and mass resolution on the results presented in this letter by applying the analysis described in Section 3 to two lower resolution re-runs of Firebox, listed in Table A1. Following the selection criteria discussed in Section 2, this results in 23 $L_*$ galaxies for the intermediate resolution volume and 25 $L_*$ galaxies for the low resolution volume. Additionally, we use 7 higher resolution Milky-Way cosmological zoom-in simulations from the Firebox-2 suite (m12s), to gain an intuition about convergence at and above the Firebox resolution.

We compare the different H\textsc{i} scale height measurements for $L_*$ galaxies at different resolutions in Figure A1. Apart from the H\textsc{i} scale height obtained from Equation 2, the scale height tends to be overestimated in the inner 3–5 kpc at lower resolutions, by at most factor of two at the galactic centre in the lowest resolution Firebox (FB256). Beyond that the measurements converge, showing excellent agreement at intermediate galactocentric radii, although the low resolution Firebox galaxies tend to underpredict disc flaring at large galactocentric radii. The higher resolution simulations follow the resolution trend seen in the Firebox. Differences between the m12s with Firebox are most pronounced when using Equations 1 and 2 to estimate the scale height, a consequence of the differences in gravitational potential, velocity dispersion and surface density radial profiles of the m12s (potentially exacerbated by the smaller sample size). Scale heights measured from the vertical volume density distribution (Gaussian fit and half-width half-maximum), are well converged for Firebox and the m12s. We conclude that our results do not have a strong resolution dependence.

This paper has been typeset from a \TeX/\LaTeX\ file prepared by the author.
Figure A1. Comparing the effect of different spatial resolutions on the measured H\textsc{i} scale heights of the Firebox $L_*$ galaxies. As in Figure 2, from left to right, top to bottom, scale heights are calculated assuming hydrostatic equilibrium, from a Gaussian fit to the vertical volume density profile, the distance between the maximum and half-maximum volume density, Equation 2 and as distance from the mid-plane enclosing 63 per cent of the total H\textsc{i} mass. Coloured, dashed lines show the median of the samples: purple for the fiducial resolution Firebox, magenta for the intermediate resolution Firebox (FB512) and orange for the low resolution Firebox (FB256). Shaded regions indicate the error on the median determined via bootstrapping. Dark blue data points show the median scale height of the 7 higher resolution cosmological-zoom Fire-2 $L_*$ galaxies (m12s). Especially at intermediate radii, the simulations show excellent convergence.