CLEO Results on Upsilons

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I present highlights of recent CLEO studies on the Υ(1S), Υ(2S), and Υ(3S). The preliminary results presented here include the first observation of a hadronic cascade in the Υ system not involving pions, Υ(3S) → γχb1,2(2P) followed by χb1,2(2P) → ωΥ(1S). I also present preliminary results on a search for inclusive ψ production in Υ(1S) decays. I end with a search for nine different two body hadronic decays of the Υ(nS), n = 1, 2, 3, resulting in the discovery of two such decays, and vastly improved upper limits on the branching ratios of the rest.

1 The CLEO Detector and CLEO Data

During the year 2002, CLEO took data at the Υ(1S), Υ(2S) and Υ(3S), resulting in approximately 20 million hadronic events on the Υ(1S), 10 million hadronic events on the Υ(2S) and 5 million hadronic events on the Υ(3S). In addition to data taken on the peaks of each resonance, data were also taken below the resonances for background purposes and in scans across the resonances.

The data used in the following studies were taken with the CLEO3 detector at the CESR electron-positron storage ring. The detector includes a silicon microvertex detector and a drift chamber, as well as a crystal calorimeter and ring imaging cerenkov (RICH) detector for hadron ID. Muon chambers surround the detector. The tracking volume is in a uniform 1.5 T magnetic field.

2 What are Upsilons?

The Υ resonances are bound states of b and b̄ quarks. The spectrum of bottomonium is very much like that seen in positronium. The Υ(4S) state is the source of B mesons at the B factories. The states in question in this talk are all below B̄B threshold. The Υ(nS) states are produced directly in e⁺e⁻ collisions. The other bottomonia states are reached by cascades from these resonances. The relative rates of cascade decays are a particularly important testing ground for Lattice QCD - should these rates be accurately predicted, our overall confidence in the use of LQCD in other arenas will be bolstered.

3 Observation of χb1,2(2P) → ωΥ(1S)

Up to now, the only transitions observed among bottomonia states have involved the emission of either photons or pion pairs. In this search we looked for the transition sequence Υ(3S) →
The search involves looking for 2 leptons from the ψ(1S) decay recoiling against a photon and \( \pi^+\pi^-\pi^0 \) in data taken at the \( \Upsilon(3S) \) resonance. Of note is that no hadron particle identification is needed for this analysis, as decays with Kaons are kinematically forbidden. Also, the mass spectrum in this sample for the dilepton mass from the \( \Upsilon(1S) \) is so clean that no lepton ID cuts are required either. The photon is required to have energy between 50 and 250 MeV. The above cuts give a very clean ω peak in the 3 pion mass spectrum. Plotting the dilepton mass against the \( \omega\gamma \) recoil mass reveals a very clean signal.

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This is the first observation of a non-pion hadronic transition in the bottomonium system.

4 Charm Production in \( \Upsilon(1S) \) Decays

The \( \Upsilon \) system offers a glue rich environment in which it is interesting to study the mechanisms of charm production. The belief is that charm should be produced through the color octet mechanism. Our study focused on searching for the inclusive production of the \( \psi \) in the decays of the \( \Upsilon(1S) \). To do this we searched for \( \psi \rightarrow e^+e^-, \mu^+\mu^- \) decays in the data taken at the \( \Upsilon(1S) \) resonance. We subtracted the luminosity and beam energy scaled contribution from continuum processes using data taken below the resonance. Additional cuts were used to suppress radiative returns to the \( \psi \) and \( \psi(2S) \). As a cross check, the same analysis method was used on \( \Upsilon(4S) \) data to verify that its results were as expected. The resulting preliminary branching fraction is found to be

\[
Br(\Upsilon(1S) \rightarrow \psi X) = (6.1 \pm 0.3 \pm 0.6) \times 10^{-4}
\]  

Also of interest is the beam energy scaled momentum spectrum of the \( \psi \)'s in this process. It appears that this spectrum is softer than would be expected from a naive color octet model, although it might be possible to address this issue with the emission of soft gluons in the theoretical calculations.

5 Two Body Hadronic Decays of the \( \Upsilon \)

To date, no purely hadronic decays of the \( \Upsilon(nS)(n = 1, 2, 3) \) have been observed. In addition, it is of interest to see if these decays in the \( \Upsilon \) system will follow what has been called variously the 14% or 12% rule seen in \( \psi \) and \( \psi' \) decay. This rule of thumb follows from the naive observation that in \( \psi(nS) \) decay, the relative production widths of the \( e^+e^- \) and \( q\bar{q} \) final states should be proportional. Extending this ansatz to the fully hadronised final states leads to the expectation that the branching ratio of \( \psi' \rightarrow XY \) should be 12% of the branching ratio of \( \psi \rightarrow XY \). This rule is grossly obeyed in the \( \psi \) system, except that the decay rate \( \psi' \rightarrow \rho \pi \) seems too small. This is occasionally referred to as the \( \rho \pi \) puzzle. In the \( \Upsilon \) system, with 3 states below B threshold, the rule translates to a 48% (72%) rule when comparing 2S (3S) to 1S decays.

This analysis concentrates on a search for \( \Upsilon \) decays to those final states that have large rates in \( \psi \) decays. The modes chosen are \( \rho \pi \), \( K^*(892)\bar{K} \), \( \rho a_2(1320) \), \( \omega f_2(1270) \), \( \phi f_2(1525) \), \( K^*(892)K_2^*(1430) \), \( b_1(1235)\pi \), \( K_1(1270)\bar{K} \), and \( K_1(1400)\bar{K} \).

Fig. 4 shows the scaled total energy for the decay products of the \( \Upsilon(1S) \) decays. Note that there are statistically significant signals for the \( \phi f_2(1525) \) and \( K_1(1400)\bar{K} \) final state, leading to our claim of first observation for these modes. All other final states lead to vastly improved...
Figure 1: Scaled Total Energy For 2 Body Hadronic Decays of the Υ(1S). The points with error bars are data from the Υ(1S), while the light histogram is the scaled expectation from non Υ(1S) data, and the dark histogram represents arbitrarily normalized signal Monte Carlo.
Table 1: Preliminary Results on Branching Ratios for \( \Upsilon \) Decays to Two Body Hadronic States - all results are in units of \( 10^{-6} \) - limits are 90% confidence level upper limits.

| Final State | \( \Upsilon(1S) \) | \( \Upsilon(2S) \) | \( \Upsilon(3S) \) |
|------------|----------------|----------------|----------------|
| \( \rho \pi \) | < 4 | < 11 | < 22 |
| \( K^*(892)K \) | < 11 | < 8 | < 14 |
| \( \rho a_2(1320) \) | < 19 | < 24 | < 30 |
| \( \omega f_2(1270) \) | < 7 | < 11 | < 8 |
| \( \phi f_2'(1525) \) | \( 7.2 \pm 1 \) | < 17 | < 14 |
| \( K^*(892)K_2^*(1430) \) | < 19 | < 32 | < 28 |
| \( b_1(1235) \pi \) | < 8 | < 12 | < 18 |
| \( K_1(1270)K \) | 8 | < 11 | < 17 |
| \( K_1(1400)K \) | \( 14.3 \pm 2 \) | < 33 | < 22 |

upper limits on the branching ratios in the range of \( 10^{-5} \) to \( 10^{-6} \), as summarized in Table 1. These limits indicate that branching fractions to these fully hadronic states are suppressed by a factor of more than 100 relative to those seen in the \( \psi \) system. Unfortunately, this means that much more data, and much more work will be needed to test the 12% rule in the \( \Upsilon \) system.

6 Conclusions

CLEO is now fully exploiting the world’s largest sample of \( \Upsilon \) decays. We have reported on the observation of \( \chi_{b1,2} \rightarrow \omega \Upsilon(1S) \), which the first observation of a transition in the \( \Upsilon \) system not involving a pion or photon. We have also measured the rate for inclusive \( \psi \) production in \( \Upsilon(1S) \) decays, and measured the momentum spectrum of the \( \psi \)'s. This will hopefully be useful in shedding light on the mechanisms of gluonic production of \( c \bar{c} \) pairs. We have also presented the first observation of two hadronic decay modes in the \( \Upsilon(1S) \) system and set stringent limits on 7 others from the \( \Upsilon(1S) \) and 9 each from the \( \Upsilon(2S) \) and \( \Upsilon(3S) \). All results reported here are preliminary.

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References

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