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Composite Topological Objects in Topological Superfluids

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Abstract—The spontaneous phase coherent precession of magnetization, discovered in 1984 by Borovik-Romanov, Bunkov, Dmitriev and Mukharskiy \cite{1} in collaboration with Fomin \cite{2}, became now an important experimental tool for study complicated topological objects in superfluid $^3$He.

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1. INTRODUCTION

Superfluid phases of $^3$He discovered in 1972 \cite{3} opened the new area of the application of topological methods to condensed matter systems. Due to the multi-component order parameter which characterizes the broken $SO(3) \times SO(3) \times U(1)$ symmetry in these phases, there are many inhomogeneous objects—textures and defects in the order parameter field—which are protected by topology and are characterized by topological quantum numbers. Among them there are quantized vortices, skyrmions and monopoles, solitons and vortex sheets, monopoles and boojums, Alice strings, Kibble–Lazarides–Shafi walls terminated by Alice strings, spin vortices with soliton tails, etc. \cite{4}. Most of them have been experimentally identified and investigated using nuclear magnetic resonance (NMR) technique, and in particular the phase coherent spin precession discovered in 1984 in $^3$He-B by Borovik-Romanov, Bunkov, Dmitriev and Mukharskiy \cite{1, 5} in collaboration with Fomin \cite{2}. Such precessing state, which has got the name Homogeneously Precessing Domain (HPD), is the spontaneously emerging steady state of precession, which preserves the phase coherence across the whole sample in the absence of energy pumping and even in an inhomogeneous external magnetic field. This spontaneous coherent precession has all the signatures of the coherent superfluid Bose–Einstein condensate of magnons (see review paper \cite{6}).

The Bose condensation of magnons in superfluid $^3$He-B had many practical applications. In Helsinki, owing to the extreme sensitivity of the Bose condensate to textural inhomogeneity, the phenomenon of Bose condensation has been applied to studies of topological defects by the HPD spectroscopy.

2. SUPERFLUID PHASES OF LIQUID $^3$He

In bulk liquid $^3$He there are two topologically different superfluid phases, $^3$He-A and $^3$He-B \cite{7}. One is the chiral superfluid $^3$He-A with topologically protected Weyl points in the quasiparticle spectrum. In the ground state of $^3$He-A the order parameter matrix has the form

$$A_{\alpha i} = \Delta_\alpha \rho e^{i \phi} \hat{d}_\alpha (\hat{e}_i^\prime + i \hat{e}_2^\prime), \quad \hat{I} = \hat{e}_1 \times \hat{e}_2,$$

where $\hat{d}$ is the unit vector of the anisotropy in the spin space due to spontaneous breaking of $SO(3)_S$ symmetry of spin rotations, and $\hat{e}_1$ and $\hat{e}_2$ are mutually orthogonal unit vectors; and $\hat{I}$ is the unit vector of the anisotropy in the orbital space due to spontaneous breaking of orbital rotations $SO(3)_L$ symmetry. The $\hat{I}$-vector also shows the direction of the orbital angular momentum of the chiral superfluid, which emerges due to spontaneous breaking of time reversal symmetry. The chirality of $^3$He-A has been probed in several experiments \cite{8–10}.

Another phase is the fully gapped time reversal invariant superfluid $^3$He-B. In the ground state of $^3$He-B the order parameter matrix has the form

$$A_{\alpha i} = \Delta_\alpha \rho e^{i \phi} R_{\alpha i},$$

where $R_{\alpha i}$ is the real matrix of rotation. $R_{\alpha i} R_{\beta j} = \delta_{\alpha \beta}$. This phase has topologically protected gapless Majorana fermions living on the surface (see reviews \cite{11, 12} on the momentum space topology in superfluid $^3$He).

In $^3$He confined in the nematicallly ordered aerogel (nafen), new phase becomes stable - the polar phase of $^3$He \cite{13, 14}, with the order parameter

$$A_{\alpha i} = \Delta_\rho e^{i \phi} \hat{d}_\alpha \hat{m}_j.$$
3. STRINGS WITH SOLITONIC TAIL

3.1. HPD and Combined Spin–Mass Vortices with Solitonic Tail

There are different types of the topological defects in the $^3$He-B. Among them there are the conventional mass vortices with the $N$ winding number of the phase $\Phi$, and the $Z_2$ spin vortex—the nontrivial winding of the matrix $R_{\Phi}$. Due to spin-orbit coupling the spin vortex serves as the termination line of the topological soliton wall. Because of the soliton tension the spin vortex moves to the wall of the vessel and escapes the observation. However, the help comes from the mass vortices. The mass and spin vortices are formed by different fields. They do not interact since they “live in different worlds”. The only instance, where the spin and mass vortices interact, arises when the cores of a spin and a mass vortex happen to get close to each other and it becomes energetically preferable for them to form a common core. Thus by trapping the spin vortex on a mass vortex the combined core energy is reduced and a composite object $Z_2$-string + soliton + mass vortex, or spin-mass vortex is formed. This object is stabilized near the edge of the vortex cluster in the rotating cryostat (see Fig. 1).

These combined objects have been observed and studied using HPD spectroscopy [18, 19]. The additional absorption observed in the homogeneously precessing domain (HPD) is proportional to the soliton area $A = lh$, where $h$ is the height of the container, and $l$ is the length of the cross-section of the soliton. In the rotating container the length $l$ is given by the width of the counterflow vortex-free zone, which is regulated by changing the angular velocity of rotation $\Omega$ at fixed number $N$ of vortices in the cell:

$$l(\Omega) = R \left(1 - \frac{\Omega l(N)}{\Omega}\right),$$

where orbital vector $\hat{n}$ is fixed by the nafen strands. The reason for the appearance of the polar phase in nafen is the analog of the Anderson theorem applied for the polar phase in the presence of the columnar defects (nafen strands), see [15, 16]. While for all the other phases of superfluid $^3$He the transition temperature is suppressed by these impurities.

The polar phase is the time reversal invariant super-fluid, which contains Dirac nodal ring in the fermionic spectrum [16, 17].

The half-quantum vortices (HQVs) were originally suggested to exist in the chiral superfluid $^3$He-A [20]. The half-quantum vortex represents the condensed matter analog of the Alice string in particle physics [21] The HQV is the vortex with fractional circulation of superfluid velocity, $N = 1/2$. It is topologically confined with the fractional spin vortex, in which $\mathbf{d}$ changes sign when circling around the vortex:

$$\mathbf{d}(r)e^{i\phi(r)} = \left(\hat{x}\cos\frac{\phi}{2} + \hat{y}\sin\frac{\phi}{2}\right)e^{i\phi},$$

when the azimuthal coordinate $\phi$ changes from 0 to $2\pi$ along the circle around this object, the vector $\mathbf{d}(r)$ changes sign and simultaneously the phase $\Phi$ changes by $\pi$, giving rise to $N = 1/2$. The order parameter (5) remains continuous along the circle. While a particle that moves around an Alice string flips its charge, the quasiparticle moving around the half-quantum vortex flips its spin quantum number. This gives rise to the Aharonov-Bohm effect for spin waves in NMR experiments [22].

However, before being experimentally observed in $^3$He-A, the HQVs were first observed in another topo-
logical phase of $^3\text{He}$—the polar phase [23]. The reason for that is that in $^3\text{He}$-A the spin-orbit interaction chooses the preferable orientation for the vector $\hat{d}$ describing the spin degrees of freedom of the order parameter. This leads to formation of a soliton interpolating between two degenerate vacua with $\hat{d} = \hat{1}$ and $\hat{d} = -\hat{1}$. The energy of soliton prevents the nucleation of the Alice strings in $^3\text{He}$-A.

In contrast, in the polar phase the spin-orbit interaction can be controlled not prohibit the formation of HQVs. In the absence of magnetic field, or if the field is along the nafen strands the spin-orbit interaction does not lead to formation of the solitons attached to the spin vortices. As a result the half-quantum vortices become energetically favorable and appear in the rotating cryostat if the sample is cooled down from the normal state under rotation.

Nevertheless the solitons help to observe the Alice string first in polar phase and after that in the polar distorted A-phase. In the polar phase, when the orientation of the magnetic field is tilted with respect to aerogel strands, the spin-orbit interaction generates the solitons between the half-quantum vortices. But the Alice strings are strongly pinned by the nafen strands, and the soliton cannot shrink (see Fig. 2). The HQVs are identified due to peculiar dependence of the NMR frequency shift on the tilting angle of magnetic field [23]. The NMR experiments also allow measure the density of the Alicie string by measuring the soliton density.

Due to the strong pinning, the Alice strings formed in the polar phase by rotation of the superfluid or by the Kibble-Zurek mechanism, survive the transition to the $^3\text{He}$-A (actually to the distorted A-phase) [24].

4. HPD AND KLS WALL BETWEEN ALICE STRINGS

4.1. Nonaxisymmetric Vortex in $^3\text{He}$-B as KLS Wall Bounded by Alice Strings

The mass vortices in $^3\text{He}$-B are presented in several forms. In particular, a pair of spin-mass vortices may form a molecule, where the soliton serves as chemical bond. As a result one obtains the doubly quantized vortex, i.e. with $N = 2$ circulation quanta (see Fig. 1). Such vortex molecules have been also identified in HPD spectroscopy [18, 19].

The “conventional” $N = 1$ vortex has also an unusual structure in $^3\text{He}$-B. Already in the first experiments with rotating $^3\text{He}$-B the first order phase transition has been observed, which has been associated with the transition inside the vortex core [25]. It was suggested that at the transition the vortex core becomes non-axisymmetric, i.e. the axial symmetry of the vortex is spontaneously broken in the vortex core [26, 27]. This was confirmed in the further experiments where the coherently precessing magnetization was used [28].

In the weak coupling BCS theory, which is applicable at low pressure, such vortex can be considered as splitted into two half-quantum vortices connected by the domain wall [29, 30], which is the analog of the Kibble–Lazarides–Shafi wall bounded by cosmic strings [31]. The separation between the half-quantum vortices increases with decreasing pressure.

The phenomenon of the additional symmetry breaking in the core of the topological defect has been also discussed for cosmic strings [32]. The spontaneous breaking of the electromagnetic $U(1)$ symmetry in the core of the cosmic string has been considered, due to which the core becomes superconducting.

For the $^3\text{He}$-B vortices, the spontaneous breaking of the $SO(2)$ symmetry in the core leads to the Goldstone bosons—the modes in which the degeneracy parameter, the axis of anisotropy $b$ of the vortex core, is oscillating. The homogeneous magnon condensate, the HPD state, has been used to study the structure and twisting dynamics of this non-axisymmetric core. The coherent precession of magnetization excites the vibrational Goldstone mode via spin-orbit interaction. Moreover, due to spin-orbit interaction the precessing magnetization rotates the core around its axis with constant angular velocity. In addition, since the core was pinned on the top and the bottom of the container, it was possible even to screw the core (see Figs. 3...
and 4). Such a twisted core corresponds to the Witten superconducting string with the electric supercurrent along the core. The rigidity of twisted core differs from that of the straight core, which is clearly seen in HPD experiments (see Fig. 4).

Oscillations of the vortex core under coherent spin precession also lead to the observed radiation of acoustic magnon modes [33].

4.2. Alice Strings with KLS Wall in Polar distorted B-Phase

In the vortices with asymmetric cores the equilibrium distance between the Alice strings is rather small. The essentially larger KLS walls between the strings have been observed in the B-phase in nafen [24] (see Fig. 5). It appeared that the Alice strings formed in the polar phase by rotation of the superfluid or by the Kibble–Zurek mechanism, survive the transition to the 3He-B (actually to the distorted B-phase). They remain pinned, in spite of the formation of the KLS walls between them.

This allows us to study the unique properties of the KLS wall. In particular, the KLS wall separates two degenerate vacua with different signs of the tetrad determinant, and thus between the “spacetime” and “antispacetime” [34].

5. COMBINED OBJECTS TO BE OBSERVED

5.1. Multi-Quantum Vortex as Closed Vortex Sheet

In the chiral superfluid, the superfluid velocity \( \mathbf{v}_s \) of the chiral condensate is determined not only by the condensate phase but also by the orbital triad \( \hat{e}_1, \hat{e}_2 \) and \( \hat{1} \):

\[
\mathbf{v}_s = \frac{\hbar}{2m} (\nabla \Phi + \hat{e}_1 \nabla \tilde{\Phi} + \hat{e}_2 \nabla \tilde{\Phi}^{*}),
\]

where \( m \) is the mass of the 3He atom. As distinct form the non-chiral superfluids, where the vorticity is presented in terms of the quantized singular vortices with the phase winding \( \Delta \Phi = 2\pi N \) around the vortex core, in 3He-A the vorticity can be continuous. The contin-

Fig. 3. The vortex in 3He-B with the non-axisymmetric core as a pair of Alice strings connected by Kibble–Lazarides–Shafi wall. The HPD with its coherent precession of magnetization is used to twist the core. The vortex with twisted core is analogous to Witten superconducting string with the electric current along the string core [32].

Fig. 4. HPD absorbtion as the function of the tilting angle \( \eta \) of magnetic field in case of the Witten strings with twisted core (filled circles) and strings with untwisted core (open circles). The estimated critical angle at which the tilted field prevents twisting by HPD is in agreement with experiment.

Fig. 5. The Alice string terminating the Kibble–Lazarides–Shafi wall in the polar distorted B-phase in nematical aerogel. Due to the pinning of Alice string by the aerogel strands the KLS wall can be arbitrarily long: the wall tension is unable to unpin the string. In addition to the KLS wall there is also the soliton tail of the string. As a result one has the triple object: KLS wall + Alice string + soliton.
uous vorticity is represented by the texture of the unit vector \( \hat{l} \) according to the Mermi-Ho relation [35]:

\[
\nabla \times \mathbf{v}_s = \frac{\hbar}{4m} e_{ijk} \nabla I_j \times \nabla I_k.
\] (7)

Experimentally the continuous vorticity is typically observed in terms of skyrmions (or the Anderson–Toulouse–Chechetkin vortices [36, 37]), see the upper part of Fig. 8. Each skyrmion has \( N = 1 \) circulation quanta. There are different scenarios in which the vortex sheets with different geometries are prepared in the experiments (see [40]).

In 1994 a new type of continuous vorticity has been observed in \(^3\)He-A—the vortex texture in the form of the vortex sheets suggested by Landau and Lifshitz for the description of the rotating superfluid [43]. (bottom) The element of the vortex sheet in \(^3\)He-A. The vortex sheet is the soliton, which contains kinks in terms of merons. Each meron has circulation quantum \( N = 1 \). There are different scenarios in which the vortex sheets with different geometries are prepared in the experiments (see [40]).

Fig. 6. (top) Typical vortex sheet in \(^3\)He-A in rotating container. It mimics the system of the equidistant cylindrical vortex sheets suggested by Landau and Lifshitz for the description of the rotating superfluid [43]. (bottom) The element of the vortex sheet in \(^3\)He-A. The vortex sheet is the soliton, which contains kinks in terms of merons. Each meron has circulation quantum \( N = 1 \). There are different scenarios in which the vortex sheets with different geometries are prepared in the experiments (see [40]).

In principle, using the vortex sheet one may construct the continuous vortices with arbitrary even number \( N = 2k \) circulation quanta. This is the soliton forming the closed cylindrical surface, which contains

\[
N = 1 \quad N = 1 \quad N = 1 \quad N = 1 \quad N = 1
\]

Fig. 7. (left) The multi-quantized vortex can be stabilized as the closed vortex sheet: cylindrical soliton with merons [41]. The tension of the soliton is compensated by repulsion of vortices (merons). (right) The cosmic analog of this composite object: cosmic necklace [53]. Monopoles and/or antimonopoles are joined together by flux tubes.

\( N \) “quarks”—merons [41, 42] (see Fig. 7 left for \( N = 6 \)). However, such multi-quantum vortices are still waiting for their observation.

5.2. Monopoles, Necklaces and Monopole Lattices with Alice Strings

Another object which is waiting for its observation in \(^3\)He-A is the vortex terminated by hedgehog [44, 45]. This is the condensed matter analog of the electroweak magnetic monopole and the other monopoles connected by strings [46]. The hedgehog-monopole, which terminates the vortex, exists in particular at the interface between \(^3\)He-A and \(^3\)He-B. The topological defects living on the surface of the condensed matter system or at the interfaces are called boojums [47]. They are classified in terms of relative homotopy groups [48]. Boojums terminate the \(^3\)He-B vortex-

In 1994 a new type of continuous vorticity has been observed in \(^3\)He-A—the vortex texture in the form of the vortex sheets [38–40], see Fig. 6 top with a single vortex sheet in container. Vortex sheet is the topological soliton with kinks, each kink representing the continuous Mermin–Ho vortex with \( N = 1 \) circulation of superfluid velocity (Fig. 6 bottom).

In principle, using the vortex sheet one may construct the continuous vortices with arbitrary even number \( N = 2k \) circulation quanta. This is the soliton forming the closed cylindrical surface, which contains

\[
N = 1 \quad N = 1 \quad N = 1 \quad N = 1 \quad N = 1
\]

Fig. 8. Skyrmion in the A-phase splits into two merons. Each meron is terminated by boojum—the point topological objects, which lives at the interface between A-phase and B-phase. Boojum also plays the role of the Nambu monopole, which terminates the string—the \( N = 1 \) vortex on the B-side of the interface.

Fig. 8. Skyrmion in the A-phase splits into two merons. Each meron is terminated by boojum—the point topological objects, which lives at the interface between A-phase and B-phase. Boojum also plays the role of the Nambu monopole, which terminates the string—the \( N = 1 \) vortex on the B-side of the interface.
strings with \( N = 1 \). The boojums do certainly exist on the surface of rotating \(^3\)He-A and at the interface between the rotating \(^3\)He-A and \(^3\)He-B [49] (see Fig. 8). However, at the moment their NMR signatures are too weak to be resolved in NMR experiments in \(^3\)He. But the vortex terminated by the hedgehog-monopole was observed in cold gases [50].

The HPD state has its own topological defects [51], and among them are the spin and orbital monopoles connected by string in Fig. 9.

In particle physics the monopoles terminating strings are called Nambu monopoles [52]. Several monopoles connected by strings may form the multi-monopole objects, such as necklace in Fig. 7 (right) [53]. This is similar to the vortex sheet necklace in Fig. 7 (left).

In \(^3\)He-A the analogs of Nambu monopoles and Alice strings may form the more complex combinations. This is because the monopole serves as a source or sink of \( N = 2 \) circulation quanta, and thus can be the termination point of 4 Alice strings with \( N = 1/2 \) each. This in particular allows construct the 2D and 3D lattices of monopoles, in Fig. 10 and in Fig. 11 correspondingly.

6. CONCLUSIONS

Here we considered several types of the topological confinement. The composite topological objects were experimentally observed in superfluid \(^3\)He by using the unique phenomenon of HPD—the spontaneously formed coherent precession of magnetization discovered by the Borovik-Romanov group in Kapitza Institute. With HPD spectroscopy, two key objects have been identified in \(^3\)He-B: spin-mass vortex [\( Z_2 \) spin vortex + soliton + mass vortex] and non-axisymmetric vortex [Alice string + Kibble–Lazarides–Shafi wall + Alice string]. One may expect the other more complicated examples of the topological confinement of the objects of different dimensions. The complicated composite objects, such as nexus, live also in the momentum space of topological materials [54].
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SPELL: 1. hedgehog, 2. preferable, 3. vacua, 4. unique, 5. evacua