Two GEM telescopes, each consisting of three 10x10 cm² triple-GEM chambers were built, tested and operated by the Hampton University group. The GEMs are read out with APV25 frontend chips and FPGA based digitizing electronics developed by INFN Rome.

The telescopes were used for the luminosity monitoring system at the OLYMPUS experiment at DESY in Germany, with positron and electron beams at 2 GeV. The GEM elements have been recycled to serve in another two applications: Three GEM elements are used to track beam particles in the MUSE experiment at PSI in Switzerland. A set of four elements has been configured as a prototype tracker for phase 1a of the DarkLight experiment at the Low-Energy Recirculator Facility (LERF) at Jefferson Lab in Newport News, USA, in a first test run in summer 2016.

The Hampton group is responsible for the DarkLight phase-I lepton tracker in preparation. Further efforts are ongoing to optimize the data acquisition speed for GEM operations in MUSE and DarkLight. An overview of the group’s GEM detector related activities will be given.
1. Introduction

The Gas Electron Multiplier (GEM) detector is a charge amplification device, which provides high-resolution particle tracking in a high particle-rate environment [1]. The gas inside the GEM detector is ionized locally by an intercepting ionizing particle, and the released electrons are guided through three layers of GEM foils representing regions with a large electric field, to undergo avalanche amplification. The avalanche produces a total charge large enough to be processed by suitable electronics. The schematics of a triple-GEM chamber is shown in Fig. 1 (left).

Two GEM telescopes were built, tested and operated by Hampton University (HU) to be utilized in the OLYMPUS experiment [2, 3, 4]. Each telescope was 70 cm long and consisted of three $10 \times 10 \text{ cm}^2$ triple-GEM chambers, which were 30-40 cm separated from each other. The GEM chambers were designed at MIT inspired by the GEM design of the COMPASS experiment at CERN [5]. An exploded view is shown in Fig. 1 (middle). The GEM foils, 5-µm copper-clad 50-µm Kapton layers perforated with a dense hole pattern, were produced by TechEtch, Inc. in Massachusetts, USA. An operating voltage of $\sim 400 \text{ V}$ is applied between the top and bottom side of the GEM foil, causing a large electric field inside the holes. One passive voltage divider (VD) per chamber is used to generate a voltage ladder for high voltage between and across the three GEM foils. Each VD was tested for a maximum voltage of 4500 V. The total operating voltage was between $-4,100 \text{ V}$ and $-4,200 \text{ V}$ between the cathode layer and the readout anode at ground. The three GEMs of one telescope are operated with pre-mixed Ar:CO$_2$ (70:30) gas at a flow rate of $\sim 1 \text{ L/h}$ in series. The amplified charge cloud is directed toward the anode, a 2D readout structure at ground potential, to collect the induced charge. The 2D readout plane is a flexible PCB consisting of copper strips in one orientation and pads connected through vias for virtual strips in crossed orientation at a pitch of 400 µm. A microscope view of the readout layer is shown in Fig. 1 (right). The design with pads and vias makes charge sharing between both coordinates more reproducible.

The GEM signals are read out with APV25 frontend chips [6], each to serialize 128 readout channels, and FPGA-based digitizing electronics developed by INFN Rome [7]. Each GEM chamber has 250 channels for each coordinate axis, read out with two APVs on either side, resulting in 500 readout channels per GEM element or a total of 1500 readout channels per telescope. The VME-based Multi-Purpose Digitizer (MPD) board, v3.0, was initially used with our GEMs to communicate with 12 APV25 front-end cards. The MPD board transmits both the control and configuration signals and digitizes the analog signals. The telescopes were triggered by $12 \times 12 \text{ cm}^2$ scintillator pads read out by silicon photomultipliers (SiPM).

![Figure 1: Left: Schematics of a triple-GEM. Middle: GEM layout (exploded view). Right: Readout layer.](image)
2. GEMs for OLYMPUS

The OLYMPUS experiment (pOsitron-proton and eLectron-proton elastic scattering to test the hYpothesis of Multi-Photon exchange USing DoriS) was carried out to study the effect of “Two Photon Exchange” [2, 3]. The experiment utilized the positron and electron beams of 2 GeV and 100 mA at the DORIS multi-GeV storage ring at DESY in Hamburg, Germany. The 12° luminosity monitoring system consisted of three triple GEMs interleaved with three Multi-Wire Proportional Chambers (MWPC) and thin scintillators in the front and back, shown in Fig 2 (left). The high resolution of the HU GEMs helped to characterize the MWPCs and to reconstruct and calibrate the vertex distributions. The HU GEMs performed at OLYMPUS with high efficiency ∼95% and high spatial resolution ∼70 µm [3, 4].

3. GEMs for MUSE

The MUSE experiment (MUon Scattering Experiment) is underway at Paul Scherrer Institute (PSI) in Switzerland to study the “Proton Radius Puzzle” [8, 9]. MUSE uses a low-energy e/π/µ beam at PSI’s πM1 beamlne for a direct test if muon-proton (µp) and electron-proton (ep) scatterings are different. Mixed beams of (e-/π-/µ-) or (e+/π+/µ+) separated by time of flight information impinge on a liquid H₂ target to simultaneously measure absolute cross sections and form factors for ep and µp scattering, allowing to extract the proton charge radius with high precision. NSF and DOE approved MUSE for funding in 2016 to run in 2018-2019.

The focused πM1 beam envelope has a width of a few cm with angular variation of ∼100 mrad at the scattering target. Event-by-event beam particle tracking is needed to reconstruct the incoming particle track for a precise scattering angle determination. The Hampton group is responsible for the beam particle tracking detector, shown in Fig. 2 (right). The existing 10×10 cm² 3-GEM telescope from OLYMPUS is utilized and has been re-configured for this purpose. Test runs have been in progress since 2013 to characterize the πM1 beam, to test detector prototypes (scintillators, Cherenkov counters, straw tubes) and to study and optimize the GEM performance.

The MUSE beam line is operated at up to 3.3 MHz beam flux with multiple beam momenta from 115-210 MeV/c. A beam flux density on the GEMs of up to 5 kHz/mm² is expected, which

![Figure 2: Left: HU GEMs together with MWPCs at OLYMPUS at DESY. Right: HU GEMs as a beam line particle tracker at MUSE at PSI.](image-url)
is still well below the COMPASS limit of 25 kHz/mm² [5]. The HU GEMs were reconfigured for MUSE and have been successfully operated under MUSE conditions.

4. GEM performance at MUSE

The APVs in OLYMPUS and MUSE have been operated in single-sample peak mode. For each acquired event, all channels were digitized by an ADC and recorded. As an initial step of the GEM analysis, baseline fluctuations were investigated at different stages. In each event, only a few channels respond to a hit with a reading above the baseline, while all other channels will show a digitized reading of the baseline itself. Each channel shows an individual average baseline value or pedestal. The baseline fluctuates by the so-called common-mode noise, a correlated up and down variation of the ADC value for all channels of one APV in every event. The common-mode fluctuation about a mean value can be assessed for any given event by averaging the raw ADC values of those channels that did not participate in a hit. At low occupancy this is the case for most channels. The common mode fluctuation about a mean value (obtained from averaging many events) is then subtracted from each channel event-by-event. The pedestal of each channel can be determined from averaging the raw ADC reading for many events, either before and after the common-mode correction. Figure 3 (left) shows the common-mode corrected pedestals, averaged over many events, for two APVs (250 channels) attached to one example GEM axis. For the same dataset, Fig. 3 (middle) shows the noise variation for the conditions of no subtractions (iliac), only common-mode correction (red), only pedestal subtraction (magenta), and common-mode corrected pedestal subtraction (blue). A residual noise level (RMS) of 40-60 channels is typically observed, depending on the cabling scheme and operating conditions.

The GEM cluster finding algorithm developed for OLYMPUS [3] is being used for the MUSE GEM data analysis presented here. The ADC readings after common-mode and pedestal subtraction for the strips of the X and Y axis are scanned for local maxima. The readings of the next two (±2) neighbored strips are summed to determine the charge of the cluster candidate. A charge threshold is applied to accept the cluster candidate. The cluster coordinates are obtained from an ADC-weighted average of the strips involved in the cluster. The charge of 1D clusters in X and Y

![Figure 3: Left: Common-mode corrected average pedestals for each strip for one GEM coordinate axis (250 channels). Middle: Noise variation for the same dataset before and after common-mode correction and bin-to-bin pedestal subtraction. Right: Identified cluster map for the US GEM.](image-url)
belonging to the same originating detector hit will be correlated. Charge sharing between the 1D clusters is imposed to determine the correct combination for a 2D cluster location if two or more 1D clusters are found. Figure 3 (right) shows the cluster map for a selected GEM (“US”). One can independently check the number of 2D clusters recorded for each event on each GEM. This quantity is called “the cluster multiplicity”. The accumulated distribution of cluster multiplicities under beam flux conditions ranging from a few hundred kHz to a few MHz peaks around unity for each GEM, with 10-20% probability for 2 or more clusters in any given event.

Straight-line tracks are formed between the cluster candidates of the first (US) and the last (DS) GEM and projected to the middle GEM (MS). The track location at the middle GEM (MS) can be calculated and compared with the cluster location of that GEM for the so-called track residuals at the MS GEM. A distance between the projected track and the detected cluster location of < 1 cm is required in order to accept the track as valid. This requirement eliminates most of the spurious tracks originating from false clusters. Using extrapolation, the track locations intercepting with any defined plane can be determined. Figure 4 (left) shows the beam spot represented by accumulated track locations at the MUSE target position, extrapolated with the GEM telescope situated upstream of the target center. Figures 4 (middle and right) show the track angles in the horizontal and vertical planes, respectively. The track resolution can be obtained from the track residuals and are dominated by multiple scattering to a few hundred μm in MUSE. In OLYMPUS, where the tracked particles were more energetic, intrinsic spatial resolutions of ≈70 μm were observed.

Straight tracks from events with one cluster in any two GEMs can be used to map out the efficiency of the third GEM, and applied to each of three GEMs. The efficiency map for each GEM is defined by taking the ratio of the number of projected tracks (determined the other GEMs), for which at least one cluster is detected (by the considered GEM) in the vicinity of the projected track, and the total number of projected candidate tracks. In order to determine the GEM efficiency maps for the entire active area of the GEMs, they need to be fully illuminated by a defocused beam. The GEM readout was triggered with a 12×12 cm² scintillator pad, which is slightly larger than the GEM active area. Figure 5 shows the efficiency maps for the three GEM elements in the MUSE GEM telescope for the data taken at one of the MUSE test beam times at PSI. The GEM efficiencies averaged over the active area were found to be 97-99%.
5. GEMs for DarkLight

The DarkLight (Detecting A Resonance Kinematically with cLectrons Incident on a Gaseous Hydrogen Target) experiment has been proposed to search for a heavy photon, $A'$ in the mass range of 10-100 MeV/c$^2$ produced in electron-proton collisions. The search includes visible decays, $A' \rightarrow e^+e^-$ and invisible decays, $A' \rightarrow X$ in in ep $\rightarrow$ ep$A'$ [10].

DarkLight was planned to be run in different phases at the Low Energy Recirculator Facility (LERF) at Jefferson Lab in the US. The phase-Ia detector tracked leptons inside the DarkLight solenoid with a set of GEM detectors, combined with segmented scintillators for triggering. The Hampton group has been responsible for the GEM lepton tracker. A fourth GEM chamber was added to reconfigure one of the existing 3-GEM telescopes from OLYMPUS. The MPD v4.0 was used with an upgraded cabling scheme to communicate with 16 APV25 front-end cards, to process 2,000 readout channels for the new 4-GEM stack. The GEM chambers were mounted in a slightly angled fixed frame, which is about 12 cm tall, displayed in Fig. 6 (left and middle). The setup was installed inside the DarkLight solenoid at LERF to detect lepton tracks in the experiment, seen in Fig. 6 (right). DarkLight Phase-Ia was carried out in 2016 aiming for LERF accelerator studies. LERF delivered a 100 MeV electron beam onto a windowless hydrogen gas target. The GEMs were operated with and without the target in place. A similar cluster finding algorithm as in MUSE can be used for the Darklight GEM data analysis. However, the solenoidal magnetic field will cause the tracks to be circular in the transverse plane, or spiraling in 3D. Phase Ib (low-energy Moller) is pursued at MIT’s HVRL, while Phase Ic (search for an $A'$ near 17 MeV/c$^2$) is in the planning stage.

6. Outlook and Conclusion

Using a different cable scheme, the residual noise at MUSE has been further improved to $\sim$30-50 ADC channels RMS. The GEM data analysis software is further refined for cluster finding, and more advanced GEM tracking algorithms are under development for higher multiplicities in a high beam flux environment. Further efforts are ongoing to optimize the performance with the MPD v4.0, and to increase the GEM data acquisition rate to $\approx$2 kHz @ 20% dead time (0.1ms/event). To this end we are in the process of upgrading the DAQ system to take advantage of advanced VME features (MBLT, 2eVME and 2eSST), which requires supporting VME hardware (64x crates and Intel controller with TSI148 chipset). In addition, we will split the VME readout into three separate crates, one for each GEM, which will increase the DAQ rate by another factor 3. Multi-

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**Figure 5:** GEM efficiency maps for all three GEMs at MUSE.
Development of GEM Detectors at Hampton University

Anusha Liyanage

Figure 6: Left: DarkLight GEM layout. Middle: DarkLight 4-GEM setup fixed in a slightly angled fixed frame. Right: GEMs installed inside the solenoid.

A sample APV readout will allow to extract event timing to eliminate randoms or background hits not originating from the triggered event. Lastly, the Hampton GEM group is responsible for the construction of (25×40 cm$^2$) GEMs for DarkLight Phase Ic presently underway.

The Hampton University GEM chambers have been in use for three major experiments, OLYMPUS at DESY in Germany, MUSE at PSI in Switzerland and DarkLight at Jefferson Lab in the US, in stable operation with > 97% efficiency and better than < 100 µm spatial resolution.

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References

[1] F. Sauli, GEM: A new concept for electron amplification in gas detectors, Nucl. Instrum. Methods Phys. Res. A 386, 531 (1997).
[2] R. Milner et al., The OLYMPUS Experiment, Nucl. Instrum. Methods Phys. Res. A 741, 1 (2014).
[3] B.S. Henderson et al., Hard Two-Photon Contribution to Elastic Lepton-Proton Scattering Determined by the OLYMPUS Experiment, Phys. Rev. Lett. 118, 092501 (2017).
[4] O. Ates, GEM Luminosity Monitors for the OLYMPUS Experiment to Determine the Effect of Two-Photon Exchange, Ph.D. thesis, Hampton University, Hampton, Virginia, 2014.
[5] M.C. Altunbas et al., Construction, test and commissioning of the triple-GEM tracking detector for COMPASS, Nucl. Instrum. Methods Phys. Res. A 490, 177 (2002).
[6] M.J. French et al., Design and results from the APV25, a deep sub-micron CMOS front-end chip for the CMS tracker, Nucl. Instrum. Methods Phys. Res. A 466, 359 (2001).
[7] V. Bellini et al., GEM tracker for high luminosity experiments at the JLab Hall A, JINST 7, C05013 (2012).
[8] R. Gilman et al., Studying the Proton “Radius” Puzzle with µp Elastic Scattering (Proposal for MUSE at PSI), arXiv:1303.2160 [nucl-ex] (2013).
[9] R. Gilman et al., Technical Design Report for the Paul Scherrer Institute Experiment R-12-01.1: Studying the Proton “Radius” Puzzle with µp Elastic Scattering, arXiv:1709.09753 [physics.ins-det] (2017).
[10] J. Balewski et al., The DarkLight Experiment: A Precision Search for New Physics at Low Energies, arXiv:1412.4717 [physics.ins-det] (2014).