Correlations of electrons from heavy flavor decay in p+p, d+Au and Au+Au collisions

Anne Sickles, for the PHENIX Collaboration Brookhaven National Laboratory

E-mail: anne@bnl.gov

In relativistic heavy ion collisions heavy flavor probes are crucial to understand the interactions between partons and the produced hot nuclear matter. Measurements in p+p collisions provide information about how the heavy quarks are produced and fragment and in d+Au collisions are sensitive to possible effects from cold nuclear matter. Azimuthal correlation measurements involving heavy flavor probes are complementary to single particle spectra measurements and provide additional information about production and interactions of heavy quarks. Measurements of electrons with heavy flavor decay with other hadrons from the event can provide information about how the heavy quark interacts with the produced matter and can be compared to similar measurements from light hadron correlations. Correlations between electrons from heavy flavor decay with muons, also from heavy flavor decay, can provide further information about heavy flavor production and cold nuclear matter effects in d+Au collisions with a very clean signal.

We present PHENIX results for electron-hadron correlations in p+p and Au+Au collisions and electron-muon correlations in p+p and d+Au collisions and discuss the implications of these measurements.
Angular correlations particles from the decay and fragmentation of heavy quarks is of great interest in hadronic and nuclear collisions. In $p+p$ collisions these measurements can aid in the understanding of the heavy flavor production mechanisms. In small nuclear systems, such as $d+Au$ collisions these correlations can be sensitive to cold nuclear matter and gluon saturation effects. In heavy ion collisions, these measurements are crucial to understanding heavy quark energy loss in hot nuclear matter. A substantial suppression of electrons from the decay of $D$ and $B$ mesons with respect to expectations from $p+p$ collisions has been observed at mid-rapidity [1]. This was unexpected if the hard partons moving through the hot nuclear matter lost energy primarily via gluon bremsstrahlung because of the charm and bottom quarks’ large mass [2]. However interpretation of correlation measurements requires measurements in $p+p$ collisions to establish a reliable baseline.

The PHENIX experiment at RHIC has measured azimuthal correlations of electrons from heavy flavor decay and charged hadrons in $p+p$ and $Au+Au$ collisions and the correlations between electrons and muons at forward rapidity, where both leptons are from heavy flavor decay in $p+p$ and $d+Au$ collisions at $\sqrt{s_{NN}}=200\text{GeV}$.

In PHENIX heavy flavor is measured via leptons from the semileptonic decay of $D$ and $B$ mesons. PHENIX has excellent electron identification capabilities at mid-rapidity $|\eta|<0.35$. At more forward and backward rapidities, $1.2<|\eta|<2.4$, muons are measured. With the planned installation of silicon vertex detectors at both mid-rapidity (2010) and forward rapidity (2011) further separation between charm and bottom decay products can be made, but currently they are combined into a single sample. For electrons at mid-rapidity PHENIX has measured the fraction of electrons from bottom decay by looking at the invariant mass of $e^{-}K$ pairs and found it to be in agreement with theoretical calculations [3]. In the electron transverse momentum, $p_{T,e}$, range of interest here, $p_{T,e}<4.5\text{ GeV}/c \approx 10\%-50\%$ of the electrons are from $B$ decay; the fraction increases with increasing $p_{T,e}$.

Azimuthal angle correlations at mid-rapidity have been used extensively at RHIC to complement single particle yield measurements. Correlations between trigger particles, in this case electrons from heavy meson decay, and charged hadrons are measured. Pairs correlated only through event wise correlations are subtracted [4]. Remaining pairs at this momentum range are correlated through being products of the same hard parton-parton scattering. Measured electron-hadron ($e_{inc}−h$) correlations are a weighted average of the correlations from electrons from heavy meson decay and those from background sources:

$$ Y_{e_{inc}−h}(\Delta \phi) = \frac{N_{eHF}Y_{eHF}−h + N_{\ell\ell}Y_{\ell\ell}−h}{N_{eHF} + N_{\ell\ell}} $$

(1)

$Y$ is the correlated yield of hadrons as a function of $\Delta \phi$ per electron. Here the background electrons are dominantly from $\pi^0$ and $\eta$ Dalitz decays and photon conversions in air and detector material. The photons which convert are also primarily from $\pi^0$ and $\eta$ decay. The background contribution can be determined from other measurements (including, for example, the $\pi^0$ $p_T$ spectrum), so the ratio $R_{HF} = \frac{N_{eHF}}{N_{\ell\ell}}$ can be determined [5, 1]. $Y_{eHF}(\Delta \phi)$ can be expressed as (leaving off the $\Delta \phi$ dependence of all $Y$ terms):

$$ Y_{eHF}−h = \frac{(R_{HF} + 1)Y_{e_{inc}−h} − Y_{\ell\ell}−h}{R_{HF}}. $$

(2)
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Figure 1: Left: Away side conditional yields for $e_{HF} - h$ correlations for four electron $p_T$ bins and eight hadron $p_T$ bins. Right: $I_{AA}$ on the away side for $2.0 < p_{T,e} < 3.0 \text{ GeV/c}$ electron triggers as a function of $p_{T,h}$.

The remaining unknowns are the correlations of the electrons from background sources with hadrons, $Y_{e_{bkg}} - h$. These are determined from measured photon-hadron correlations (like the background electrons, inclusive photons are dominantly from $\pi^0$ and $\eta$ decay at these $p_T$s); simulations are used to account parent meson distributions in the decay electron and decay $\gamma$ samples [6]. Fig. 1(a) shows the azimuthally integrated away side yield ($\Delta \phi \approx \pi$), as a function of the $p_T$ of the associated hadron for four electron $p_T$ selections.

Fig. 1(b) shows the away side yield in Au+Au divided by the yield in p+p collisions, $I_{AA}$. The result suggests some away side suppression of the parton opposing the leading heavy quark traverses the hot nuclear matter and is qualitatively consistent with the $I_{AA}$ measured in hadron-hadron correlations [7].

In heavy quark fragmentation most of the quark momentum is carried by the heavy meson [8, 9], thus the near side $e_{HF} - h$ is expected to be dominated by pairs in which both particles come from the heavy meson decay. Due to the decay kinematics, we then expect the Gaussian widths of the near side to be wider for $e_{HF} - h$ than $e_{inc} - h$. The widths from both categories of electrons are shown in Fig. 3. The measured near side widths for $e_{HF} - h$ are wider than $e_{phot} - h$ and are in agreement with widths from charm production in PYTHIA [10].

Electron-muon correlations were proposed long ago as a clean $c\bar{c}$ signal in heavy ion collisions [11]. One advantage to measuring electron-muon correlations is that both particles are from the decay of heavy flavor. The correlated charm signal produces electrons and muons of opposite sign from each other, thus most backgrounds such as jets and combinatoric pairs are removed by subtracting like sign pairs. Remaining backgrounds are due to leptons not from charm decay. Fig. 3(a) shows the final electron-muon azimuthal correlations in p+p collisions [12]. Electrons have $p_T > 0.5 \text{ GeV/c}$ and $|\eta| < 0.35$ and muons have $p_T > 1.0 \text{ GeV/c}$ and $1.4 < |\eta| < 2.1$. The near side peak is missing because the pairs are from back-to-back $c\bar{c}$ pairs; any near side correlation would not extend over the large rapidity difference between the two leptons. The same measurement in d+Au collisions with the muon going in the deuteron going direction ($1.4 < \eta < 2.1$) is shown in Fig. 3(b). Obviously, a d+Au event has a larger average pair yield than a p+p event. Thus the d+Au data has been scaled down by the number of average number binary nucleon-nucleon
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**Figure 2:** Gaussian widths as a function of $p_{T,h}$ for four electron $p_T$ bins. Circles show the $e_{HF} - h$ widths and open diamonds show $e_{phot} - h$ widths. Pluses show $e_{HF} - h$ widths from PYTHIA and the triangles show an earlier PHENIX measurement of $e_{HF} - h$ widths with lower statistics and $1.0 < p_T < 4.0$ GeV/c.

collisions in the event sample, $N_{coll}$. This is appropriate because the creation of the $c\bar{c}$ pairs is a hard process which should scale with $N_{coll}$. At small $\Delta \phi$ the scaled $d+Au$ and $p+p$ collisions are consistent with each other, but at $\Delta \phi \approx \pi$ the $d+Au$ pair yields are significantly suppressed with respect to $p+p$ collisions.

With the muon at forward rapidity (deuteron going direction) these correlations are sensitive to the low $x$ region of the Au nucleus. Study of the $x$ region of the Au nucleus probed in this kinematic region shows some sensitivity to $x \leq 0.01$ where saturation effects could possibly be significant [13].

In summary, PHENIX has measured the azimuthal correlations of leptons from heavy flavor decay in three collision systems. These measurements are sensitive to the details of heavy flavor production and its modification in nuclear collisions. A significant suppression of rapidity separated electron-muon pairs is observed in $d+Au$ collisions. Further studies are ongoing. In Au+Au collisions we see suppression of away side associated hadrons similar to those from hadron triggered correlations. While the uncertainties in these measurements are large, measurements will be improved in the near future from upgrades to the PHENIX detector, including silicon vertex detectors at mid-rapidity and forward rapidity and increased statistics.

**References**

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Figure 3: Azimuthal distributions of opposite sign electron-muon pairs per event after background subtraction in p+p (left) and d+Au (right) collisions. The d+Au distribution has been scaled down by the mean number of binary collisions in the event sample in order to be compared to p+p collisions.

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