Perturbation of Riemann-Hilbert jump contours: smooth parametric dependence with application to semiclassical focusing NLS

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Abstract

A perturbation of a class of scalar Riemann-Hilbert problems (RHPs) with the jump contour as a finite union of oriented simple arcs in the complex plane and the jump function with a $z \log z$ type singularity on the jump contour is considered. The jump function and the jump contour are assumed to depend on a vector of external parameters $\vec{\beta}$. We prove that if the RHP has a solution at some value $\vec{\beta}_0$ then the solution of the RHP is uniquely defined in a some neighborhood of $\vec{\beta}_0$ and is smooth in $\vec{\beta}$. This result is applied to the case of semiclassical focusing NLS.

1 Introduction

We study a type of scalar Riemann-Hilbert problem that appears in semiclassical or small dispersion calculations of integrable systems. We prove the smooth dependence of the solution on a crucial parameter. Our direct motivation comes from the semiclassical focusing NLS equation [2, 14, 15, 16]

$$i\varepsilon \partial_t q + \varepsilon^2 \partial^2_x q + 2|q|^2 q = 0$$

with the initial condition

$$q(x, 0) = A(x)e^{ixS(x)}, \quad A(x) > 0, \quad \mu \geq 0$$

in the limit as $\varepsilon \to 0$. For any value of $\varepsilon > 0$, the solution process requires solving a $2 \times 2$ matrix Riemann-Hilbert problem (RHP) in the complex plane of the spectral parameter $z$.

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1 the equation as written here corrects a mistake in the numerical coefficients of equation (1.1) of [20].
of an underlying Lax pair operator [20, 21]. The quantities $x$, $t$, and $\mu$ appear as parameters. The Riemann-Hilbert approach is relevant to other integrable systems in general as established by [17] and it is a major tool in the asymptotic analysis of integrable systems as established by the discovery of the steepest descent method [8, 9]. The asymptotic methods also apply to orthogonal polynomial asymptotics [4, 5], and through [12, 13], to random matrices [1, 6, 10, 11].

When $\varepsilon \to 0$, the steepest descent method [8, 9], with the implementation of the $g$-function mechanism [7, 20], reduces the asymptotic calculation to a $2 \times 2$ "model" RHP which is exactly solvable. The reduction occurs through a series of transformations of the original RHP. These are facilitated by factorizations of the RHP jump-matrix and appropriate RHP contour deformations. As a result, the main contributions to the solution are isolated. Higher order contributions can be calculated iteratively. Characteristic of nonlinearity, the main contributions are linked to arcs in the complex plane that form the "contributing" part of the contour of the model RHP. The process is in the spirit of the steepest descent method for the asymptotic evaluation of integrals, where, however, the contributions are localized near points.

In the nonlinear problem, the arc endpoints are crucial quantities that control the shape of the nonlinear waveforms that develop in the small $O(\varepsilon)$ spatiotemporal scale. The number and the position of these endpoints in the complex plane vary with $x$ and $t$ in the large $O(1)$ scale, thus, modulating the waveform in space-time. They are often referred to as the modulation parameters or as branchpoints, as they arise naturally as the branchpoints of a square root.

The $g$-function mechanism reduces this process to solving a scalar RHP and, thus, determining the branchpoints and all other data needed for the model problem. Conceptually, the scalar RHP problem identifies how crucial features of the eigenfunctions and of the potential vary in the large space-time scale and facilitates the rigorous derivation of the solution of the matrix RHP. It is, thus, an analogous entity to the eikonal equation of linear PDEs. The function $g$ that solves the scalar RHP problem is, essentially, the same phase function that appears in the direct semiclassical scattering problem. In the RHP it is considered as a function of the spectral variable, at fixed $x$ and $t$. The eikonal equation it addresses it as a function of $x$ and $t$ at a fixed value of the spectral variable.

The smooth perturbation of the solution of the scalar RHP with respect to parameters $x$ and $t$ was obtained in [18], excluding a crucial parameter $\mu$ at which the input function $f(z)$ to the scalar RHP had a $z \log z$ singularity. The following formula for the branch points (see below for the definition of $\alpha_j$ and other quantities in the formula) was derived

$$\frac{\partial \alpha_j(\vec{\beta})}{\partial \beta_k} = -\frac{2\pi i}{D(\vec{\alpha}, \vec{\beta})} \frac{\partial K(\alpha_j, \vec{\alpha}, \vec{\beta})}{\partial \beta_k} \frac{f'(\zeta, \vec{\beta})}{(\zeta - \alpha_j(\vec{\beta})) R(\zeta, \vec{\alpha})} d\zeta, \quad \beta_2 = x, \quad \beta_3 = t. \quad (3)$$

We prove that the formula holds for the parameter $\beta_1 = \mu$ as well and the dependence is smooth, meaning that the contour, the jump matrix, and the solution of the scalar RHP evolve smoothly. The significance of this is that this and the continuation method will allow extending long-time estimates of the position of branchpoints to parameter regions beyond the ones for which such estimates are known. In the NLS case this means extending to regions of $\mu$ in which solitons appear.
The Perturbation theorem 4.1 has an immediate application to NLS with initial data being a semiclassical approximation \[20\] of

\[ q(x, 0) = -\text{sech} \left( x \right) e^{-\frac{i}{\epsilon} \int_0^x \tanh(s) ds}, \quad \mu \geq 0. \]  

This family of initial conditions is interesting because of a transition at \( \mu = 2 \) with solitonless interval \( (\mu \geq 2) \) and soliton interval \( 0 < \mu < 2 \). It has been studied in a number of papers \[20, 21, 22\] The semiclassical limit has been completely analyzed for \( \mu \geq 2 \) while for the soliton case \( 0 < \mu < 2 \) the answer is known only for some finite times. The RH approach runs into difficulties with the error estimates and was not able to continue past a certain curve. Numerical experiments have shown absence of any noticeable transition in the behavior of end points \( \alpha_j(\mu) \) at the critical value \( \mu = 2 \) \[3\]. The Perturbation theorem for NLS 5.3 establishes this fact rigorously.

## 2 The scalar Riemann-Hilbert problem (RHP): background

The following scalar RHP is studied in \[15\]. In this section we provide an overview. The contour of RHP, labeled \( \gamma \) is an unknown of the problem, except for the fact that it is a connected, finite, non self intersecting, typically open curve, that is partitioned into \( 2N + 1 \) arcs \( [\alpha_0, \alpha_1], [\alpha_1, \alpha_2], [\alpha_2, \alpha_3], \ldots \) by the points \( \alpha_0, \alpha_1, \ldots, \alpha_{2N+1} \), that have the natural ordering of their indices \( 0, 1, 2, \ldots, 2N + 1 \) along the oriented contour. The number \( N = 0, 1, 2, \ldots \) is given in the problem. A scalar function \( g(z) \) is sought that is analytic and bounded in \( \mathbb{C} \setminus \gamma \). The jump condition on the \( N + 1 \) main arcs \( [\alpha_0, \alpha_1], [\alpha_2, \alpha_3], \ldots, [\alpha_{2N}, \alpha_{2N+1}] \) (they include the arcs at the two ends of the contour) is

\[ g_+ + g_- = f + W \]  

where \( g_\pm(z) \) are the nontangential limiting values of \( g(z) \) from the positive/negative side of the contour, \( f(z) \) is a given function, that constitutes the main input to the problem, and
$W$ is a real constant to be determined. The constant $W$ generally takes different values on different arcs. The contour is sought in the domain of analyticity of $f(z)$, except for a finite number of its points where $f(z)$ can be non-analytic [15]. In the case of NLS, $f(z)$ must be determined from the initial data [2].

The condition on the remaining $N$ complementary arcs, that interlace with the main arcs is

$$g_+ - g_- = \Omega$$

(6)

where $\Omega$ is again a real constant to be determined and takes different values on different arcs.

A bounded function $g(z)$ satisfying these conditions can be written explicitly, with the aid of the radical

$$R(z) = \left( \prod_{j=0}^{2N+1} (z - \alpha_j) \right)^{\frac{1}{2}}$$

where

$$\lim_{z \to \infty} \frac{R(z)}{z^{N+1}} = -1,$$

(7)

in which the main arcs are the branchcuts. To write the formula for $g$ we introduce some notation. The main arcs are labeled $\gamma_{m,j} = (\alpha_{2j}, \alpha_{2j+1}), j = 0, 1, ..., N$. The complementary arcs are labeled $\gamma_{c,j} = (\alpha_{2j-1}, \alpha_{2j}), j = 1, 2, ..., N$. The subscript + in $R(\zeta)_+$ indicates that the value of $R$ on the branchcut (that is, on the contour) is taken as the limiting value from the left side of the contour. The following $g(z)$ (see [20]) is analytic and bounded in $\mathbb{C} \setminus \gamma$ and satisfies the jump conditions [5] and [6]

$$g(z) = \frac{R(z)}{2\pi i} \left[ \oint_{\gamma} \frac{f(\zeta)}{(\zeta - z)R(\zeta)_+} d\zeta + \sum_{j=0}^{N} \oint_{\gamma_{m,j}} \frac{W_j}{(\zeta - z)R(\zeta)_+} d\zeta + \sum_{j=1}^{N} \oint_{\gamma_{c,j}} \frac{\Omega_j}{(\zeta - z)R(\zeta)_+} d\zeta \right].$$

(8)

4
The real constants $W_j$ and $Ω_j$ are chosen so that $g(z)$ is $O(1)$ as $z → ∞$, and, hence, analytic at $z = ∞$. They are calculated as the unique solution to a linear system of moment conditions that guarantee this (see equations (14) of paper [18]).

Without loss of generality, we can fix one of the constants $W_j$, say $W_0 = 0$. If $W_0 ≠ 0$, this can be archived by recalibrating the RH problem by considering $g(z) − ½ W_0$ and subtracting $W_0$ to $g(∞)$ and all $W_j$ on main arcs $γ_{m,j}$.

Rather than to perform integration along arcs we prefer to integrate along loops surrounding these arcs. Thus, we define the function $h(z)$

$$h(z) = \frac{R(z)}{2πi} \left[ \oint_{γ} \frac{f(ζ)}{(ζ - z)R(ζ)} dζ + \sum_{j=0}^{N} \oint_{γ_{m,j}} \frac{W_j}{(ζ - z)R(ζ)} dζ + \sum_{j=1}^{N} \oint_{γ_{c,j}} \frac{Ω_j}{(ζ - z)R(ζ)} dζ \right], \tag{9}$$

in which the oriented loop $γ$ surrounds contour $γ$, $γ_{c,j}$ encircles the complementary arc $γ_{c,j}$ (notice different orientations of the two halves of the loop), and $γ_{m,j}$ encircles the main arc $γ_{m,j}$. All loops are oriented clockwise, $z$ lies inside the contour $γ$ and outside the loops $γ_{m,j}$ and $γ_{c,j}$.

Passing from arcs to loops introduces factor of 2 and $z$ cutting through the deforming loop around $γ$ introduces a residue in the first integral.

Thus the relation of $g(z)$ and $h(z)$ is,

$$g(z) = \frac{1}{2}(h(z) + f(z)). \tag{10}$$

The jump conditions for $h$ along the contour are simpler than the jumps of $g(z)$, thus, $h(z)$ is the most natural object in the Riemann-Hilbert analysis. Results are easily reformulated in terms $g(z)$, whose main advantage is its analyticity off the contour. The analyticity of $h$ of the contour, is compromised at the non-analytic points of $f$ and on the real axis.

In order to evaluate the limits of function $h(z)$ at a branchpoint, coming from the left or from the right side of the RHP contour, we allow $z$ to cross inside the loops of the main and complementary arcs adjacent to the branchpoint and we correct by introducing the corresponding residue. Denoting the expression in the bracket for this positioning of $z$ by

$$B(z) = \oint_{γ} \frac{f(ζ)}{(ζ - z)R(ζ)} dζ + \sum_{j=0}^{N} \oint_{γ_{m,j}} \frac{W_j}{(ζ - z)R(ζ)} dζ + \sum_{j=1}^{N} \oint_{γ_{c,j}} \frac{Ω_j}{(ζ - z)R(ζ)} dζ, \tag{11}$$

we obtain, for $B$ calculated with $z$ inside the two loops,

$$h(z) = W_j ± Ω_j + \frac{R(z)B(z)}{2πi} \text{ or } h(z) = W_j ± Ω_{j+1} + \frac{R(z)B(z)}{2πi}, \tag{12}$$

where the first equation applies near the branchpoint $α_{2j}$, the second equation applies near the branchpoint $α_{2j+1}$, and $+$ or $−$ applies when $z$ is to the left or right of the contour respectively. The $±$ term is zero at the endpoints of the contour $α_0$ and $α_{2N+1}$.

It is important to notice that $B(z)$ is analytic at $z = α_j$ and can thus be expanded to a convergent power series in the neighborhood of $α_j$,

$$B(z) = ν_{1,j} + ν_{2,j}(z - α_j) + ν_{3,j}(z - α_j)^2 + \ldots, \tag{13}$$
with coefficients that are analytic with respect to all of the branchpoints.

The growth/decay properties of \( e^{\pm ih(z)/\varepsilon} \) near the RHP contour \( \gamma \) is a crucial to the semiclassical analysis of matrix RHP. The requirement that the constants \( W_i \) and \( \Omega_j \) be real is a fall-out of this. An additional fall-out is the following sign conditions for \( h \):

\[
\Im h = 0 \text{ on main arcs}; \quad \Im h < 0 \text{ left and right of main arcs}
\]

(14)

\[
\Im h > 0 \text{ on complementary arcs}
\]

(15)

The last condition must also apply on two "extension arcs", or arcs extending the contour to \(+\infty\) on one side and to \(-\infty\) on the other.

For the sign conditions to be satisfied, it is necessary, but not sufficient that

\[
B(\alpha_j) = 0, \quad j = 0, 1, 2, \ldots, 2N + 1,
\]

equivalently,

\[
\nu_{1,j}(\alpha_j) = 0, \quad j = 0, 1, 2, \ldots, 2N + 1.
\]

(17)

equivalently,

\[
h(z) = \text{analytic} + O((z - \alpha_j)^3), \quad \text{as } z \to \alpha_j, \quad j = 0, 1, 2, \ldots, 2N + 1.
\]

(18)

equivalently,

\[
g(z) = \text{analytic} + O((z - \alpha_j)^3), \quad \text{as } z \to \alpha_j, \quad j = 0, 1, 2, \ldots, 2N + 1.
\]

(19)

equivalently,

\[
h'(\alpha_j) = 0, \quad j = 0, 1, 2, \ldots, 2N + 1.
\]

(20)

The necessity of the coefficient of the linear term of \( B \) being zero is intimately related to a simple fact about \( \sqrt{z} \), with the positive semiaxis as a branchcut playing the role of a main arc. With a sign determination that gives \( \Im \sqrt{z} < 0 \) on both sides of the branchcut, as required by the sign condition, \( \Im \sqrt{z} \) never turns positive, making the connection to a complementary arc impossible. Figure 3 describes how the sign conditions can be satisfied with \( \sqrt{z} \).

In [18], \( B(z) \) is expressed as a ratio of two determinants through,

\[
B(z) = \frac{2\pi i}{D} K(z),
\]

(21)

where

\[
K(z) = \frac{1}{2\pi i} \begin{vmatrix}
    f_{\gamma_m,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,1} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_m,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,1} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_N,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_N,1} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_N,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_N,1} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_c,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,1} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_c,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,1} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_m,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,N} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_m,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,N} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_c,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,N} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_c,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,N} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_m,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,1} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_m,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,1} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_c,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,1} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_c,1} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,1} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_m,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,N} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_m,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_m,N} \frac{z^{N-1}d\xi}{R(\xi)} \\
    \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
    f_{\gamma_c,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,N} \frac{z^{N-1}d\xi}{R(\xi)} & \ldots & f_{\gamma_c,N} \frac{d\xi}{R(\xi)} & \ldots & f_{\gamma_c,N} \frac{z^{N-1}d\xi}{R(\xi)}
\end{vmatrix}
\]
Figure 3: Sign structure of $\Im \sqrt{z}$ near the origin where the branchcut is chosen along the positive semiaxis.

and

$$D = \begin{vmatrix} \frac{f_{\gamma_1,1}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_1,1}}{R(\zeta)} & \cdots & \frac{f_{\gamma_m,1}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_m,1}}{R(\zeta)} \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\ \frac{f_{\gamma_1,N}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_1,N}}{R(\zeta)} & \cdots & \frac{f_{\gamma_m,N}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_m,N}}{R(\zeta)} \\ \frac{f_{\gamma_1,1}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_1,1}}{R(\zeta)} & \cdots & \frac{f_{\gamma_m,1}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_m,1}}{R(\zeta)} \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\ \frac{f_{\gamma_1,N}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_1,N}}{R(\zeta)} & \cdots & \frac{f_{\gamma_m,N}}{R(\zeta)} & \cdots & \frac{\zeta^{N-1} f_{\gamma_m,N}}{R(\zeta)} \end{vmatrix}.$$  \tag{23}

and the following relations are shown

$$g(z) = \frac{R(z)}{2D} K(z), \quad \tag{24}$$

where $z$ lies outside of $\gamma$ and outside all $\gamma_{c,j}$ and $\gamma_{m,j}$ and

$$h(z) = \frac{R(z)}{D} K(z), \quad \tag{25}$$

where $z$ lies inside of $\gamma$ and outside all $\gamma_{c,j}$ and $\gamma_{m,j}$.

For the rest of the paper we consider function $K(z)$ which is a constant multiple of $B(z)$. From [16] the modulation equations to determine the branchpoints $\alpha_j$ are the following

$$K(\alpha_j) = 0, \quad j = 0, \ldots, 2N + 1. \quad \tag{26}$$
3 Smooth parametric dependence of the branchpoints \( \alpha_j \)

By external parameters \( \vec{\beta} = (\beta_1, \beta_2, \ldots) \) we think either \( \vec{\beta} \in \mathbb{R}^m \) or \( \vec{\beta} \in \mathbb{C}^m, \ m \geq 1. \)

**Definition 3.1.**

Let a simple contour \( \gamma_0 \) consist of a finite union of finite length oriented simple arcs in the complex plane \( \gamma_0 = (\cup_{m,j} \gamma_{m,j}) \cup (\cup_{c,j} \gamma_{c,j}) \) with the distinct end points \( \alpha_0 = \{ \alpha_j \}_{j=0}^{2N+1} \), as shown in Fig. [1]. For a fixed vector of external parameters \( \vec{\beta}_0 \), we define

\[
\mathbb{L}(\gamma_0, \alpha_0, \vec{\beta}_0)
\]

to be the set of functions \( f(z, \vec{\beta}) \) defined in some open neighborhood of \( (\gamma_0, \vec{\beta}_0) \) which have the form

\[
f(z, \vec{\beta}) = c(\vec{\beta}) (z - z_0) \log (z - z_0) + A(z, \vec{\beta}), \quad (27)
\]

with \( z_0 \in \gamma_{m,j} \cup \cup_k \{ \alpha_k \} \) for some \( j \) (without loss of generality, let \( z_0 \in \gamma_{m,0} \setminus \{ \alpha_0, \alpha_1 \} \)); and where any branch of the logarithm is chosen so that the branchcut does not intersect \( \gamma_0 \).

Assume \( c(\vec{\beta}) \) and \( A(z, \vec{\beta}) \) are twice continuously differentiable in parameters \( \vec{\beta} \) in some open neighborhood of \( \vec{\beta}_0 \); and for each \( \vec{\beta} \) in some open neighborhood of \( \vec{\beta}_0 \), \( A(z, \vec{\beta}) \) is analytic in \( z \) in some open neighborhood of \( \gamma_0 \).

Let \( z_0 \), the point of logarithmic singularity of \( f \), depend on one of the parameters, say \( \beta_1 \). Assume \( z_0 = z_0(\beta_1) \) is continuously differentiable in \( \beta_1 \). For simplicity, assume that \( z_0(\beta_1 + \Delta \beta_1) \) approaches to \( \gamma_0 \) as \( \Delta \beta_1 \to 0 \) non-tangentially.

Note that for fixed \( \alpha_0 \) and \( \vec{\beta}_0 \) the contour \( \gamma_0 \) aside from passing through \( \alpha_0 \) and \( z_0 \) is free to deform continuously as long as \( f \in \mathbb{L}(\gamma_0, \alpha_0, \vec{\beta}_0) \). This condition fixes the (distinct) end points of \( \gamma_0 \) and the singularity \( z_0 \). Thus a contour \( \gamma_0 \) from the definition 3.1 depends on the end points \( \alpha_0 \) and on \( \beta_1^0 \) through \( z_0, \gamma(\alpha_0, \beta_1^0) = \gamma_0 \).

**Lemma 3.2.** Let \( \gamma_0 \) be a simple oriented contour with the distinct end points \( \alpha_0 \) and assume \( f \in \mathbb{L}(\gamma_0, \alpha_0, \vec{\beta}_0) \). Then there exist open neighborhoods of \( \alpha_0 \) and \( \vec{\beta}_0 \) such that

\[
f \in \mathbb{L}(\gamma, \alpha, \vec{\beta})
\]

for all \( \alpha \) in the neighborhood of \( \alpha_0 \), for all \( \vec{\beta} = (\beta_1, \ldots) \) in the neighborhood of \( \vec{\beta}_0 \), and some contour \( \gamma = \gamma(\alpha, \vec{\beta}_1). \)

**Proof.** From the definition \( f \in \mathbb{L}(\gamma_0, \alpha_0, \vec{\beta}_0) \) implies that \( f \in \mathbb{L}(\gamma_0, \alpha_0, \vec{\beta}) \) for all \( \vec{\beta} \) in some open neighborhood of \( \vec{\beta}_0 \).

Now fix some \( \vec{\beta} \) in the neighborhood of \( \vec{\beta}_0 \) with \( f \in \mathbb{L}(\gamma_0, \alpha_0, \vec{\beta}) \). Since the form of \( f \) (27) is valid in some neighborhood of \( \gamma_0 \), then there is a neighborhood of \( \alpha_0 \) where for all \( \alpha \) the contour \( \gamma_0 \) with the end points \( \alpha_0 \) can be deformed into \( \gamma \) with the end points \( \alpha \). For example, by continuously connecting the end points \( \alpha_0 \) with the points \( \alpha \). Thus \( f \in \mathbb{L}(\gamma, \alpha, \vec{\beta}). \)

The size of the neighborhoods of \( \vec{\beta}_0 \) and \( \alpha_0 \) is determined by the distance \( \alpha_j \)’s with respect to each other and the distance between \( \gamma \) and the singularities of \( f(z) \) (other than \( z_0 \)).
By considering the loop contours \( \hat{\gamma}, \hat{\gamma}_m, \hat{\gamma}_c \) (see Figure 2) the explicit dependence of the contours of integration on the end points \( \vec{\alpha} \) is removed (for example in (35-38)). So even though \( \gamma = \gamma(\vec{\alpha}, \beta_1) \), in all our evaluations below \( \hat{\gamma} = \hat{\gamma}(\beta_1) \).

The main difficulty is the dependence of \( f(z) \) (thus the RHP (3)) and the modulation equations (26) on parameter \( \beta_1 \) which also controls the point \( z_0 \) on the contour \( \hat{\gamma} \). We show that the dependence on \( \beta_1 \) (moreover on \( \vec{\beta} \)) is smooth.

Remark 3.3. In [18] was considered the case when the contour \( \gamma \) was independent of external parameters \( \vec{\beta} \) which led to simpler conditions on \( f \) being continuous on \( \gamma \) except at finitely many points. If \( f \in \mathbb{L}(\gamma, \vec{\alpha}, \vec{\beta}) \), then all results below for \( \beta_2, \beta_3, \ldots \) follow from [18]. However the harder case of the dependence on \( \beta_1 \) is new.

Consider the system of modulation equations as a function of both \( \vec{\alpha} \) and \( \vec{\beta} \)

\[
K(\alpha_j) = K(\alpha_j, \vec{\alpha}, \vec{\beta}) = 0, \quad j = 0, 1, \ldots, 2N + 1.
\]  

(28)

Note, that the function \( K(z, \vec{\alpha}, \vec{\beta}) \) is defined by (22) and through (25) defines \( h(z, \vec{\alpha}, \vec{\beta}) \), which satisfies a scalar Riemann-Hilbert problem on \( \gamma = \gamma(\vec{\alpha}, \beta_1) \):

\[
\begin{cases}
  h_+(z) + h_-(z) = 2W_j(\vec{\alpha}, \vec{\beta}), & \text{on } \gamma_m, \quad j = 0, 1, \ldots, N, \\
  h_+(z) - h_-(z) = 2\Omega_j(\vec{\alpha}, \vec{\beta}), & \text{on } \gamma_c, \quad j = 1, \ldots, N, \\
  h(z) + f(z) \text{ is analytic in } \mathbb{C} \setminus \gamma.
\end{cases}
\]  

(29)

So every time we use \( K \) we understand that there is an underlying scalar Riemann-Hilbert problem which depends on \( \vec{\alpha} \) and \( \vec{\beta} \).

Following [18], we differentiate (28) with respect to \( \beta_k \)

\[
\sum_l \frac{\partial K(\alpha_j)}{\partial \alpha_l} \frac{\partial \alpha_l}{\partial \beta_k} + \frac{\partial K(\alpha_j)}{\partial \beta_k} = 0,
\]  

(30)

where the matrix \( \left\{ \frac{\partial K(\alpha_j)}{\partial \alpha_l} \right\}_{j,l} \) is diagonal [18] so

\[
\frac{\partial K(\alpha_j)}{\partial \alpha_j} \frac{\partial \alpha_j}{\partial \beta_k} = -\frac{\partial K(\alpha_j)}{\partial \beta_k}.
\]  

(31)

Since

\[
\frac{\partial K(\alpha_j)}{\partial \alpha_j} = \frac{D(\vec{\alpha}, \vec{\beta})}{2\pi i} \oint_{\hat{\gamma}(\beta_1)} \frac{f'(\zeta, \vec{\beta})}{(\zeta - \alpha_j)R(\zeta, \vec{\alpha})} d\zeta
\]  

(32)

we arrive to the evolution equations for \( \alpha_j \):

\[
\frac{\partial \alpha_j}{\partial \beta_k} = -\frac{2\pi i \frac{\partial K(\alpha_j)}{\partial \beta_k}}{D(\vec{\alpha}, \vec{\beta}) \oint_{\hat{\gamma}(\beta_1)} \frac{f'(\zeta, \vec{\beta})}{(\zeta - \alpha_j)R(\zeta, \vec{\alpha})} d\zeta}.
\]  

(33)

Since \( D \neq 0 \) for distinct \( \alpha_j \)'s [18], next we need to estimate the partial derivatives \( \frac{\partial K(\alpha_j)}{\partial \beta_k} \).
Lemma 3.4.

Let \( f \in \mathbb{L}(\gamma, \bar{\alpha}, \bar{\beta}_0) \), where the contour \( \gamma = \gamma(\bar{\alpha}, \beta_1) \) has fixed end points \( \bar{\alpha} \). Then there is an open neighborhood of \( \bar{\beta}_0 \) such that for all \( \bar{\beta} \) in the neighborhood of \( \bar{\beta}_0 \)

\[
\frac{\partial K(\alpha_j)}{\partial \beta_k} = \frac{1}{2\pi i} \frac{\partial}{\partial \beta_k} \left| \begin{array}{cccc}
\frac{d}{R(\zeta)} f_{\gamma_{m,1}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{m,1}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{m,1}} \\
\ldots & \ldots & \ldots & \ldots & \ldots \\
\frac{d}{R(\zeta)} f_{\gamma_{m,N}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{m,N}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{m,N}} \\
\frac{d}{R(\zeta)} f_{\gamma_{c,1}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{c,1}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{c,1}} \\
\ldots & \ldots & \ldots & \ldots & \ldots \\
\frac{d}{R(\zeta)} f_{\gamma_{c,N}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{c,N}} & \ldots & \frac{d}{R(\zeta)} f_{\gamma_{c,N}} \\
\end{array} \right|.
\]

To proceed we need the following technical lemma.

**Proof.** The idea of the proof is to consider the finite differences and take the limit. The main difficulty is the case \( k = 1 \) when both the integrand and the contour depend on \( \beta_1 \).

Denote the integral on the left in (35) as \( I_1 \)

\[
I_1(\bar{\beta}) = \oint_{\hat{\gamma}(\beta_1)} \frac{\zeta^n f(\zeta, \bar{\beta})d\zeta}{R(\zeta, \bar{\alpha})}.
\]

First, consider \( k \geq 2 \). Without loss of generality, assume that the branchcut of logarithm in \( f(z, \bar{\beta}) \) is chosen from \( z_0(\beta_1) \) horizontally (along the straight line \( z = i\Im z_0 \)) to the right near \( z_0 \) (see Fig. 4).

The contour of integration \( \hat{\gamma} \) near \( z_0 \) is pushed to the logarithmic branchcut of \( f \). Fix complex points \( \delta_1 \) and \( \delta_2 \) near \( z_0 \) with \( \Im \delta_1 = \Im \delta_2 = \Im z_0 \) (see Fig. 4). Then \( \hat{\gamma} \) is split into
two parts: $[\delta_1, \delta_2]$ and its complement. Across the logarithmic branchcut $[z_0, \delta_2]$, $f(z, \bar{\beta})$ has a jump $2\pi i (z_0 - z) c(\bar{\beta})$ and no jump across $[\delta_1, z_0]$.

Take a small $\Delta \beta_k$ and consider finite differences

$$I_1(\bar{\beta} + \Delta \bar{\beta}_k) - I_1(\bar{\beta}) = \int_{\gamma(\beta_1)} \frac{\zeta^n f(\zeta, \bar{\beta} + \Delta \bar{\beta}_k) - f(\zeta, \bar{\beta})}{\Delta \bar{\beta}_k} R(\zeta, \alpha) d\zeta,$$

where $\Delta \bar{\beta}_k = \Delta \beta_k \bar{e}_k$ with $\bar{e}_k$ to be the standard basis $(\bar{e}_k)_j = \begin{cases} 1, & j = k, \\ 0, & j \neq k. \end{cases}$

The limit $\Delta \beta_k \to 0$ can be interchanged with the integral in (40) since the integrand is uniformly bounded in $\beta_k$ on $\gamma$, $k \geq 2$.

In the case $k = 1$, there are two logarithmic branchcuts of $f(z, \bar{\beta})$ and $f(z, \bar{\beta} + \Delta \bar{\beta}_1)$. By assumption $z_0$ is not tangential to $\gamma$ and so for small $\Delta \beta_1$, $z_0(\beta_1 + \Delta \beta_1)$ is not on $\gamma$. Assume that $z_0(\beta_1 + \Delta \beta_1)$ is on the left (positive) side of $\gamma_{m,0}(\beta_1)$. The case when $z_0(\beta_1 + \Delta \beta_1)$ is on the other side of $\gamma_{m,0}(\beta_1)$ is done similarly.

Without loss of generality, assume that the branchcut of logarithm in $f(z, \bar{\beta})$ is chosen as before and in $f(z, \bar{\beta} + \Delta \bar{\beta}_1)$ the branchcut is chosen from $z_0(\beta_1 + \Delta \beta_1)$ to $z_0(\beta_1)$ and horizontally to the right (see Fig. 4). Then similarly, we choose fixed points $\delta_1$ and $\delta_2$ so that $3\delta_1 = 3z_0(\beta_1 + \Delta \beta_1)$ and $3\delta_2 = 3z_0(\beta_1)$. The contour of integration $\gamma$ is pushed to the logarithmic branchcut near $z_0$ and split into $[\delta_1, \delta_2]$ and its complement, where by $[\delta_1, \delta_2]$ we understand $[\delta_1, z_0(\beta_1 + \Delta \beta_1)] \cup [z_0(\beta_1 + \Delta \beta_1), z_0(\beta_1)] \cup [z_0(\beta_1), \delta_2]$.

Across $[z_0(\beta_1), \delta_2]$, $f(z, \bar{\beta})$ has a jump $2\pi i (z_0(\beta_1) - z) c(\bar{\beta})$. Similarly, $f(z, \bar{\beta} + \Delta \bar{\beta}_1)$ has a jump $2\pi i (z_0(\beta_1 + \Delta \beta_1) - z) c(\bar{\beta} + \Delta \bar{\beta}_1)$ on $[z_0(\beta_1 + \Delta \beta_1), z_0(\beta_1)] \cup [z_0(\beta_1), \delta_2]$.

For $k = 1$,

$$I_1(\bar{\beta} + \Delta \bar{\beta}_1) - I_1(\bar{\beta}) = \int_{\gamma(\beta_1) \setminus [\delta_1, \delta_2]} \frac{\zeta^n f(\zeta, \bar{\beta} + \Delta \bar{\beta}_1) - f(\zeta, \bar{\beta})}{\Delta \bar{\beta}_1} R(\zeta, \alpha) d\zeta$$

and

$$+ \int_{[z_0(\beta_1), \delta_2]} \frac{\zeta^n 2\pi i (z_0(\beta_1 + \Delta \beta_1) - \zeta) c(\bar{\beta} + \Delta \bar{\beta}_1) - 2\pi i (z_0(\beta_1) - \zeta) c(\bar{\beta})}{\Delta \bar{\beta}_1} R(\zeta, \alpha) d\zeta$$

Figure 4: Deforming the contours of integration for $I_1$ near $z_0$. The dashed line is the logarithmic branchcut of $\log(z - z_0)$ in $f$ and $\frac{\partial f(\zeta, \bar{\beta})}{\partial \beta_k}$ in the integrals (35, 36).
The last integral is $O(\Delta \beta_1)$ which is observed by the change of variables $y = \zeta - z_0(\beta_1 + \Delta \beta_1)$

$$\int_{[z_0(\beta_1 + \Delta \beta_1), z_0(\beta_1)]} \frac{\zeta^n 2\pi i(z_0(\beta_1 + \Delta \beta_1) - \zeta) c(\beta + \Delta \beta_1)}{R(\zeta, \bar{\alpha})} d\zeta. \quad (43)$$

The second integral (36) is done similarly. The rest of the integrals (37)-(38) do not depend on $\beta_k$ since the only dependence on $\beta$ sits in $z_0(\beta_1) \in \gamma_{m,0}$.

Thus for fixed $\bar{\alpha}$ the integrals (37)-(38) are independent of $\beta_k$.

Thus from (34) using Lemma 3.4 for $k \geq 1$ we get

\[
\begin{aligned}
\frac{\partial K(\alpha_j)}{\partial \beta_k}(\bar{\alpha}, \bar{\beta}) &= \frac{1}{2\pi i} \begin{vmatrix}
\tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) \\
\tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) \\
\tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) \\
\tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) \\
\vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\
\tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) & \tilde{f}_{\gamma,1}(\beta) \\
\end{vmatrix}
\end{aligned}
\]
and for $k + m \geq 3$

$$\frac{\partial^2 K(\alpha_j)}{\partial \beta_k \partial \beta_m} (\vec{\alpha}, \vec{\beta}) = \frac{1}{2\pi i}$$

\[
\begin{pmatrix}
\frac{f_{j,m,1}}{R(\zeta)} & \ldots & \frac{f_{j,m,1} \zeta^{N-1} \frac{d\zeta}{R(\zeta)}}{R(\zeta)} & \ldots & \frac{f_{j,m,1}}{R(\zeta)} \\
\ldots & \ldots & \ldots & \ldots & \ldots \\
\frac{f_{j,m,N}}{R(\zeta)} & \ldots & \frac{f_{j,m,N} \zeta^{N-1} \frac{d\zeta}{R(\zeta)}}{R(\zeta)} & \ldots & \frac{f_{j,m,N}}{R(\zeta)} \\
\frac{f_{c,1}}{R(\zeta)} & \ldots & \frac{f_{c,1} \zeta^{N-1} \frac{d\zeta}{R(\zeta)}}{R(\zeta)} & \ldots & \frac{f_{c,1}}{R(\zeta)} \\
\ldots & \ldots & \ldots & \ldots & \ldots \\
\frac{f_{c,N}}{R(\zeta)} & \ldots & \frac{f_{c,N} \zeta^{N-1} \frac{d\zeta}{R(\zeta)}}{R(\zeta)} & \ldots & \frac{f_{c,N}}{R(\zeta)} \\
\end{pmatrix}
\]

Define $\tilde{K}(\vec{\alpha}, \vec{\beta})$ as

$$\left\{ \tilde{K}(\vec{\alpha}, \vec{\beta}) \right\}_j = K(\alpha_j, \vec{\alpha}, \vec{\beta}), \quad j = 0, 1, \ldots, 2N + 1.$$

**Lemma 3.5.**

Let $f \in \mathbb{L}(\gamma_0, \vec{\alpha}_0, \vec{\beta}_0)$, where the contour $\gamma_0$ has the end points $\vec{\alpha}_0$. Then

$$K_j(\vec{\alpha}, \vec{\beta}) := K(\alpha_j, \vec{\alpha}, \vec{\beta}), \quad j = 0, 1, \ldots, 2N + 1,$$

is continuously differentiable in $\vec{\alpha}$ and in $\vec{\beta} = (\beta_1, \ldots)$ in some open neighborhoods of $\vec{\alpha}_0$ and $\vec{\beta}_0$ respectively, and for some contour $\gamma = \gamma(\vec{\alpha}, \vec{\beta}_1)$.

**Proof.** $\tilde{K}(\vec{\alpha}, \vec{\beta})$ is analytic in $\vec{\alpha}$ by the determinant structure and the integral entries [22], where explicit dependence on $\vec{\alpha}$ is only in the $R(z, \vec{\alpha})$ term which is analytic away from $z = \alpha_j$.

By Lemma 3.4 the partial derivatives $\frac{\partial K_j(\vec{\alpha}, \vec{\beta})}{\partial \beta_k}$ exist and by twice differentiability of $f(z, \vec{\beta})$ in $(\beta_2, \beta_3, \ldots)$ which leads to continuous $\frac{\partial^2 K(\alpha_j)}{\partial \beta_k \partial \beta_m}$, thus $\tilde{K}(\vec{\alpha}, \vec{\beta})$ is continuously differentiable in $(\beta_2, \beta_3, \ldots)$.

For $\beta_1$, the integrals in the last row of $\frac{\partial K(\alpha_j)}{\partial \beta_1}$ in (48) involve the function

$$f_{\beta_1}(z, \vec{\beta}) = c_{\beta_1}(\vec{\beta})(z-z_0(\vec{\beta})) \log(z-z_0(\vec{\beta})) + c(\vec{\beta})z'_0(\vec{\beta})(z-z_0(\vec{\beta})) - c(\vec{\beta})z'_0(\vec{\beta}) + A_{\beta_1}(z, \vec{\beta}),$$

which is integrable near $z_0$ and hence $\frac{\partial K(\alpha_j)}{\partial \beta_1}$ is continuous in $\beta_1$.

To conclude joint smoothness of $K_j$ in $\vec{\beta}$ notice that $K_j$ can be split into the integrals of the analytic and the singular logarithmic parts of $f$ and notice that the singular part of $f$, namely, $(z - z_0(\beta_1)) \log(z - z_0(\beta_1))$ only depends on $\beta_1$ as a function of $\vec{\beta}$. This allows to conclude the joint continuous differentiability of $K_j(\vec{\alpha}, \vec{\beta})$ in $\vec{\beta}$.

Thus by Lemma 3.5 the modulation equations [28]

$$K_j(\vec{\alpha}, \vec{\beta}) = K(\alpha_j) = 0$$

(52)
are smooth in $\bar{\alpha}$ and in the external parameters $\bar{\beta}$. Next we want to solve this system for $\bar{\alpha} = \bar{\alpha}(\bar{\beta})$ and conclude smoothness in $\bar{\beta}$.

For the next lemma we need $K'(z, \bar{\alpha}, \bar{\beta}) = \frac{dK}{dz}(z, \bar{\alpha}, \bar{\beta})$

$$K'(z, \bar{\alpha}, \bar{\beta}) = \frac{1}{2\pi i} \begin{vmatrix}
    f_{j_{m,1}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{m,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{m,N}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{m,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{m,N}} \frac{d\zeta}{R(\zeta)} \\
    \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\
    f_{j_{c,1}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,N}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,N}} \frac{d\zeta}{R(\zeta)} \\
    \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\
    f_{j_{c,1}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,N}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,N}} \frac{d\zeta}{R(\zeta)} \\
    \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\
    f_{j_{c,1}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,N}} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,1}} \zeta^{N-1} \frac{d\zeta}{R(\zeta)} & \cdots & f_{j_{c,N}} \frac{d\zeta}{R(\zeta)} \\
\end{vmatrix},$$

where $z$ is inside of $\gamma(\beta_1)$ and inside of $\gamma_{m,j}$ and $\gamma_{c,j}$ or $\gamma_{c,j+1}$.

**Lemma 3.6.**

*Let $f \in L(\gamma_0, \bar{\alpha}_0, \bar{\beta}_0)$, where $\bar{\alpha}_0$ and $\bar{\beta}_0$ satisfy

$$\bar{K}(\bar{\alpha}_0, \bar{\beta}_0) = \bar{0}.$$*

*Assume that for $\bar{\alpha}_0 = \{\alpha_j^0\}_{j=0}^{2N+1}$, $\lim_{z \to \alpha_j^0} K'(z, \bar{\alpha}_0, \bar{\beta}_0) \neq 0$, $j = 0, 1, \ldots, 2N + 1$.

*Then the modulation equations

$$\bar{K}(\bar{\alpha}, \bar{\beta}) = \bar{0}$$

*can be uniquely solved for $\bar{\alpha}(\bar{\beta})$ which is continuously differentiable for all $\bar{\beta}$ in some open neighborhood of $\bar{\beta}_0$ and $\bar{\alpha}(\bar{\beta}_0) = \bar{\alpha}_0$.*

**Proof.** $\bar{K}$ is continuously differentiable in $\bar{\alpha}$ and $\bar{K}$ is a continuously differentiable in $\bar{\beta}$ by Lemma 3.5.

As it was shown in [18], the matrix

$$\begin{aligned}
\left\{ \frac{\partial \bar{K}}{\partial \bar{\alpha}} \right\}_{j,l} &= \left\{ \frac{\partial K(\alpha_j)}{\partial \alpha_l} \right\}_{j,l} \\
\end{aligned}$$

(54)

is diagonal and

$$\frac{\partial K(\alpha_j)}{\partial \alpha_j} = \frac{3}{2} D \lim_{z \to \alpha_j^0} \left( \frac{h(z)}{R(z)} \right)' = \frac{3}{2} \lim_{z \to \alpha_j^0} K'(z, \bar{\alpha}, \bar{\beta}) \neq 0.$$ (55)

So

$$\det \left| \frac{\partial \bar{K}}{\partial \bar{\alpha}(\bar{\alpha}_0)} \right| = \prod_{j} \frac{\partial K(\alpha_j)}{\partial \alpha_j} \neq 0$$ (56)

under the assumptions. By the Implicit function theorem, $\bar{\alpha}(\bar{\beta})$ are uniquely defined in some neighborhood of $\bar{\beta}_0$ and smooth in $\bar{\beta}$. Note $\bar{\alpha}(\bar{\beta}_0) = \bar{\alpha}_0$ by assumption. □
Remark 3.7.
The condition \( \lim_{z \to \alpha_0^j} K'(z, \alpha_0, \beta_0) \neq 0, \ j = 0, 1, \ldots, 2N + 1 \) in Lemma 3.6 is equivalent to \( \lim_{z \to \alpha_0^j} \frac{K'(z, \alpha_0, \beta_0)}{R(z, \alpha_0)} \neq 0, \ j = 0, 1, \ldots, 2N + 1 \).

To summarize this section, given the set of \( \alpha_0, \beta_0 = (\beta_0^0, \ldots), \gamma_0 = \gamma(\alpha_0, \beta_0^0), \) and \( f \in \mathbb{L}(\gamma_0, \alpha_0, \beta_0) \), defines the function \( K(z, \alpha_0, \beta_0) \) by formula (22) which through (25) in turn defines \( h(z, \alpha_0, \beta_0) \), which satisfies a scalar Riemann-Hilbert problem on \( \gamma_0 \):

\[
\begin{cases}
  h_+(z) + h_-(z) = 2W_j(\alpha_0, \beta_0), & \text{on } \gamma_{m,j}, \ j = 0, 1, \ldots, N, \\
  h_+(z) - h_-(z) = 2\Omega_j(\alpha_0, \beta_0), & \text{on } \gamma_{c,j}, \ j = 1, \ldots, N, \\
  h(z) + f(z) \text{ is analytic in } \mathbb{C} \setminus \gamma_0.
\end{cases}
\]

(57)

We denote this scalar Riemann-Hilbert problem as \( RHP(\gamma_0, \alpha_0, \beta_0, f) \).

Lemma 3.6 states that if \( \alpha_0 \) and \( \beta_0 \) satisfy \( \overline{K}(\alpha_0, \beta_0) = 0 \), and \( RHP(\gamma_0, \alpha_0, \beta_0, f) \) has a solution with \( \lim_{z \to \alpha_0^j} K'(z, \alpha_0, \beta_0) \neq 0 \) then \( RHP(\gamma, \alpha, \beta, f) \) has a unique smooth solution for all \( \alpha = \alpha(\beta) \) and \( \beta \) in some open neighborhood of \( \beta_0 \), where \( \alpha \) and \( \beta \) solve the modulation equations \( \overline{K}(\alpha, \beta) = 0 \).

4 Perturbation theorem

Theorem 4.1. (Perturbation Theorem)
Consider a simple contour \( \gamma_0 \) consisting of a finite union of oriented simple arcs \( \gamma_0 = (\bigcup \gamma_{m,j}) \cup (\bigcup \gamma_{c,j}) \) with the distinct end points \( \alpha_0 \) and depending on parameters \( \beta_0 \) (see Figure 1). Assume \( \alpha_0 \) and \( \beta_0 \) satisfy a system of equations

\[
\overline{K}(\alpha_0, \beta_0) = 0,
\]

and let \( f \in \mathbb{L}(\gamma, \alpha_0, \beta_0) \). Let \( \gamma = \gamma(\alpha, \beta_1) \) be the contour of a RH problem which seeks a function \( h(z) = h(z, \alpha, \beta) \) which satisfies the following conditions

\[
\begin{cases}
  h_+(z) + h_-(z) = 2W_j(\alpha, \beta), & \text{on } \gamma_{m,j}, \ j = 0, 1, \ldots, N, \\
  h_+(z) - h_-(z) = 2\Omega_j(\alpha, \beta), & \text{on } \gamma_{c,j}, \ j = 1, \ldots, N, \\
  h(z) + f(z) \text{ is analytic in } \mathbb{C} \setminus \gamma,
\end{cases}
\]

(58)

where \( \Omega_j = \Omega_j(\alpha, \beta) \) and \( W_j = W_j(\alpha, \beta) \) are real constants whose numerical values will be determined from the RH conditions. Assume that there is a function \( h(z, \alpha_0, \beta_0) \) which satisfies (58) and suppose \( \frac{h'(z, \alpha_0, \beta_0)}{R(z, \alpha_0)} \neq 0 \) for all \( z \) on \( \gamma_0 \).

Then the solutions \( \alpha = \alpha(\beta) \) of the system

\[
\overline{K}(\alpha, \beta) = 0
\]

(59)

and \( h(z, \alpha(\beta), \beta) \) which solves (58) are uniquely defined and smooth in some open neighborhood of \( \beta_0 \). Moreover, \( \Omega_j(\beta) = \Omega_j(\alpha(\beta), \beta) \), and \( W_j(\beta) = W_j(\alpha(\beta), \beta) \) are defined and smooth in \( \beta \) in some open neighborhood of \( \beta_0 \).
Furthermore, for \( k \geq 1 \):

\[
\frac{\partial \alpha_j}{\partial \beta_k}(\bar{\beta}) = -\frac{2\pi i}{D(\bar{\alpha}(\bar{\beta}), \bar{\beta})} \int_{\gamma(\bar{\beta})} f_{\gamma}(\eta, \bar{\beta}) d\eta, \quad j = 1, 2, \ldots, 2N + 1, \quad (60)
\]

\[
\frac{\partial h}{\partial \beta_k}(z, \bar{\alpha}, \bar{\beta}) = \frac{R(z, \bar{\alpha})}{2\pi i} \int_{\gamma(\bar{\beta})} \frac{\partial f}{\partial \beta_k}(\xi, \bar{\beta}) d\xi, \quad (61)
\]

where \( z \) is inside of \( \gamma(\bar{\beta}) \),

\[
\frac{\partial \Omega_j}{\partial \beta_k}(\bar{\beta}) = -\frac{1}{D} \left| \begin{array}{cccc}
\frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots & \frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots \\
\frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots & \frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots \\
\vdots & \ddots & \vdots & \ddots \\
\frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots & \frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots \\
\end{array} \right| \quad (62)
\]

\[
\frac{\partial W_j}{\partial \beta_k}(\bar{\beta}) = -\frac{1}{D} \left| \begin{array}{cccc}
\frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots & \frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots \\
\frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots & \frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots \\
\vdots & \ddots & \vdots & \ddots \\
\frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots & \frac{f_{\gamma}(\xi, \bar{\beta})}{R(\zeta)} & \cdots \\
\end{array} \right| \quad (63)
\]

where \( R(\xi) = R(\xi, \bar{\alpha}(\bar{\beta})) \).

**Remark 4.2.** Formula (60) is the same as (56) in [18]. We prove that the formula is still valid when the jump contour \( \gamma \) depends on an external parameter \( (\beta_1 \text{ in our case}) \). The contour depends on \( \beta_1 \) through the logarithmic singularity \( z_0(\beta_1) \) on \( \gamma \) (see the paragraph after the definition 3.1).
Proof. By Lemma 3.6, $\alpha_j(\beta)$ are continuously differentiable in $\beta$. Formula for $\frac{\partial g}{\partial \beta_k}$ were computed above (63).

Next we compute $\frac{\partial g(z, \bar{\alpha}, \bar{\beta})}{\partial \beta_k}$ which satisfies the RHP

$$g_{\beta_k}(z) = g_{\beta_k}(z), \quad z \in \gamma_{m,j}, \quad j = 0, 1, ..., N.$$  \hspace{1cm} (64)

Then

$$\frac{\partial g}{\partial \beta_k}(z, \bar{\alpha}, \bar{\beta}) = \frac{R(z, \bar{\alpha})}{2\pi i} \int_{\gamma} \frac{\partial f}{\partial \beta_k}(\xi, \bar{\beta})}{R(\xi, \bar{\alpha})} d\xi$$  \hspace{1cm} (65)

where $z$ is outside of $\gamma$, $\frac{\partial f}{\partial \beta_k}(z)$ for $k = 1$ behaves like $\log(z - z_0)$ near $z_0$, and for $k \geq 2$ behaves like $(z - z_0)\log(z - z_0)$ near $z_0$. So $\frac{\partial h}{\partial \beta_k}(z, \bar{\alpha}, \bar{\beta})$ satisfies (61) where $z$ is inside of $\gamma$.

Constants $W_j$ and $\Omega_j$ are found from the linear system (18)

$$\int_{\gamma(\beta_1)} \frac{\xi^n f(\xi, \bar{\beta})}{R(\xi, \bar{\alpha})} d\xi + \sum_{j=1}^{N} \int_{\gamma_{c,j}} \frac{\xi^n \Omega_j}{R(\xi, \bar{\alpha})} d\xi = 0, \quad n = 0, ..., N - 1.$$  \hspace{1cm} (66)

Differentiating in $\beta_k$ and using Lemma 3.4 leads to

$$\int_{\gamma(\beta_1)} \frac{\xi^n f_{\beta_k}(\xi, \bar{\beta})}{R(\xi, \bar{\alpha})} d\xi + \sum_{j=1}^{N} \int_{\gamma_{c,j}} \frac{\xi^n \Omega_j f_{\beta_k}}{R(\xi, \bar{\alpha})} d\xi \quad + \sum_{j=1}^{N} \int_{\gamma_{m,j}(\beta_1)} \frac{\xi^n W_j f_{\beta_k}}{R(\xi, \bar{\alpha})} d\xi = 0, \quad n = 0, ..., N - 1.$$  \hspace{1cm} (67)

and since by assumption $W_j$ and $\Omega_j$ are real and they satisfy the complex conjugate system as well

$$\int_{\gamma(\beta_1)} \frac{\xi^n f_{\beta_k}(\xi, \bar{\beta})}{R(\xi, \bar{\alpha})} d\xi + \sum_{j=1}^{N} \int_{\gamma_{c,j}} \frac{\xi^n \Omega_j f_{\beta_k}}{R(\xi, \bar{\alpha})} d\xi \quad + \sum_{j=1}^{N} \int_{\gamma_{m,j}(\beta_1)} \frac{\xi^n W_j f_{\beta_k}}{R(\xi, \bar{\alpha})} d\xi = 0, \quad n = 0, ..., N - 1,$$  \hspace{1cm} (68)

or in matrix form

\[
\begin{pmatrix}
\int_{\gamma(\beta_1)} \frac{\xi^n f_{\beta_k}(\xi, \bar{\beta})}{R(\xi, \bar{\alpha})} & \sum_{j=1}^{N} \int_{\gamma_{c,j}} \frac{\xi^n \Omega_j f_{\beta_k}}{R(\xi, \bar{\alpha})} & \sum_{j=1}^{N} \int_{\gamma_{m,j}(\beta_1)} \frac{\xi^n W_j f_{\beta_k}}{R(\xi, \bar{\alpha})}
\end{pmatrix}
\begin{pmatrix}
\frac{\partial W}{\partial \beta_k} \\
\frac{\partial \Omega_j}{\partial \beta_k}
\end{pmatrix}
= 0.
\]

So $\frac{\partial \Omega_j}{\partial \beta_k}$ and $\frac{\partial W_j}{\partial \beta_k}$ satisfy (62) and (63). Note that $D \neq 0$ for distinct $\alpha_j$’s (18).

To apply this theorem to the semiclassical focusing NLS we make a slight modification of the class of the allowed functions in the definition 3.1. In the case of NLS, the real axis is an additional contour of discontinuity of $f$ because of the Schwarz symmetry condition. This additional jump contour of $f$ affects only a small neighborhood of $z_0$. 

\[17\]
Definition 4.3. We say that \( f \in \tilde{L}(\gamma_0, \vec{\alpha}_0, \vec{\beta}_0) \) if \( f \in L(\gamma_0, \vec{\alpha}_0, \vec{\beta}_0) \) with the condition (27) is replaced with

\[
\begin{align*}
  f(z, \vec{\beta}) &= c(\vec{\beta}) (z - z_0) \log (z - z_0) + (z - z_0)B(z, \vec{\beta}) + A(z, \vec{\beta}), \\
  \lim_{z \to z_0, z \in \gamma_0} f(z, \vec{\beta}) &\text{ exists and finite,}
\end{align*}
\]

(70)

where for all \( \vec{\beta} \) in some open neighborhood of \( \vec{\beta}_0 \), \( B(z, \vec{\beta}) \) has a constant jump in \( z \) near \( z_0 \) on a simple contour \( \gamma_B \) intersecting \( \gamma_0 \) only at \( z_0 \) and coinciding with the branchcut of \( \log (z - z_0) \) near \( z_0 \). Assume \( c(\vec{\beta}) \) and \( A(z, \vec{\beta}) \) satisfy the same conditions as before, and \( B(z, \vec{\beta}) \) is twice continuously differentiable in parameters \( \vec{\beta} \) in some open neighborhood of \( \vec{\beta}_0 \); and for each \( \vec{\beta} \) in some open neighborhood of \( \vec{\beta}_0 \), \( B(z, \vec{\beta}) \) is piecewise analytic in \( z \) in some open neighborhood of \( \gamma_0\{z_0\} \) in \( C\setminus \gamma_B \).

This technical modification changes the jump of \( f \) on \([\delta_1, \delta_2]\) near \( z_0 \) in the proof of Lemma 3.4, however the jump of \( f \) remains linear and by the second condition in (70) the jump of \( f \) is 0 at \( z_0 \). This automatically is correct for \( f \in L(\gamma_0, \vec{\alpha}_0, \vec{\beta}_0) \) because of the \( z \log z \) type singularity at \( z_0 \). Thus for \( f \in \tilde{L}(\gamma, \vec{\alpha}_0, \vec{\beta}_0) \), (44-47) are correct and all the above results including Lemma 3.2, Lemma 3.6 and Theorem 4.1 remain valid.

This definition was not used from the beginning since the main difficulty is to deal with the jump of logarithm in (70) while the jump of \( B \) is a technical issue which would make the exposition less clear.

5 \( \mu \)-dependence in the semiclassical focusing NLS

The function \( f(z) \) comes from a semiclassical approximation of the family of initial conditions for NLS (41) as [20]:

\[
f(z, \mu, x, t) = \left( \frac{\mu}{2} - z \right) \left[ \frac{\pi i}{2} + \ln \left( \frac{\mu}{2} - z \right) \right] + \frac{z + T}{2} \ln (z + T) + \frac{z - T}{2} \ln (z - T)
- T \tanh^{-1} \frac{T}{\mu/2} - xz - 2tz^2 + \frac{\mu}{2} \ln 2, \quad \text{when } \Im z \geq 0
\]

(71)
and

\[ f(z) = \overline{f(\overline{z})}, \quad \text{when} \quad \Im z < 0, \]  

(72)

where the branchcuts are chosen as the following: from \( \frac{\mu}{2} \) along the real axis to \(+\infty\), from \( T \) to 0 and along the real axis to \(+\infty\), from \(-T \) to 0 and along the real axis to \(-\infty\), where

\[ T = \sqrt{\frac{\mu^2}{4} - 1}, \quad \Im T \geq 0. \]  

(73)

For \( \mu \geq 2 \), \( T \geq 0 \) is real and for \( 0 < \mu < 2 \), \( T \) is purely imaginary with \( \Im T > 0 \).

Then for \( \Im z \geq 0 \)

\[ \frac{\partial f}{\partial \mu}(z,\mu) = \frac{\pi i}{4} + \frac{1}{2} \ln \left( \frac{\mu}{2} - z \right) + \ln 2 + \frac{\mu}{8T} \left[ \ln(z + T) - \ln(z - T) - 2 \tanh^{-1} \frac{2T}{\mu} \right] \]  

(74)

where \( \tanh^{-1} x = x + O(x^3), \ x \to 0 \) then for \( z \neq 0 \)

\[ \frac{\partial f}{\partial \mu}(z,\mu) = \frac{\pi i}{4} + \frac{1}{2} \ln \left( \frac{\mu}{2} - z \right) + \ln 2 + \frac{\mu}{4z} - \frac{1}{2} + O(T), \ T \to 0. \]  

(75)

So \( \mu = 2 \) is a removable singularity for \( f_\mu \) as a function of \( \mu \)

\[ \lim_{\mu \to 2, T \to 0} \frac{\partial f}{\partial \mu}(z,\mu) = \frac{\pi i}{4} + \frac{1}{2} \ln (1 - z) + \ln 2 + \frac{1}{2z} - \frac{1}{2}, \]  

(76)

which is analytic for \( z \neq 0 \) and \( z \neq 1 \).

**Lemma 5.1.** \( f(z,\mu,x,t) \) and \( f'(z,\mu,x,t) \) are analytic in \( \mu \) for \( \mu > 0 \), \( x > 0 \), \( t > 0 \), for all \( z, \Im z \neq 0 \).

**Proof.** Consider

\[ f'(z,\mu,x,t) = -\frac{\pi i}{2} - \ln \left( \frac{\mu}{2} - z \right) + \frac{1}{2} \ln \left( z^2 - \frac{\mu^2}{4} + 1 \right) - x - 4tz, \]  

(77)

which analytic in \( \mu > 0 \), for \( \Im z \neq 0 \).

For \( \mu > 0, \mu \neq 2, f(z,\mu) \) is clearly analytic in \( \mu \) for \( \Im z \neq 0 \). At \( \mu = 2 \) \( (T = 0) \) we find the power series of \( f(z,\mu,x,t) \) in \( T \) and show that it contains only even powers. Since

\[ T^{2k} = \left( \frac{\mu^2}{4} - 1 \right)^k = \frac{(\mu + 2)^k}{4^k} (\mu - 2)^k \]  

(78)

it will show analyticity of \( f(z,\mu,x,t) \) in \( \mu \).

Start with expanding in series at \( T = 0 \)

\[ \frac{1}{\mu/2} = \sqrt{1 + T^2}^{-1} = \sum_{k=0}^{\infty} c_k T^{2k}, \quad \ln (z \pm T) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \left( \pm \frac{T}{z} \right)^n. \]  

(79)

Then the terms in (71) are analytic in \( \mu \)

\[ \frac{z + T}{2} \ln (z + T) + \frac{z - T}{2} \ln (z - T) \]  

(80)
\[ z \ln z - z \sum_{n \text{ is even}} \frac{1}{n} \left(\frac{T}{z}\right)^n + T \sum_{n \text{ is odd}} \frac{1}{n} \left(\frac{T}{z}\right)^n \tag{81} \]

\[ = z \ln z + \sum_{k=1}^{\infty} \frac{1}{2k(2k-1)} z^{2k-1} T^{2k}, \tag{82} \]

which has only even powers of \( T \) and is analytic in \( \mu \) for \( \Im z \neq 0 \). Consider the inverse hyperbolic tangent term in (71) and take into account that \( \tanh^{-1} z \) is an odd function

\[ T \tanh^{-1} \frac{T}{\mu/2} = T \tanh^{-1} \frac{T}{\sqrt{1+T^2}} = T \tanh^{-1} \sum_{k=0}^{\infty} c_k T^{2k+1} \tag{83} \]

\[ = T \sum_{k=0}^{\infty} \tilde{c}_k T^{2k+1} = \sum_{k=0}^{\infty} \tilde{c}_k T^{2k+2}, \tag{84} \]

which also has only even powers of \( T \).

So

\[ f(z) = \left( \frac{\mu}{2} - z \right) \left[ \frac{\pi i}{2} + \ln \left( \frac{\mu}{2} - z \right) \right] + \frac{z + T}{2} \ln (z + T) + \frac{z - T}{2} \ln (z - T) \]

\[- T \tanh^{-1} \frac{T}{\mu/2} - xz - 2tz^2 + \frac{\mu}{2} \ln 2 \tag{85} \]

is analytic in \( \mu \) for \( \mu > 0, x > 0, t > 0, \Im z \neq 0 \). \( \square \)

**Lemma 5.2.**

Let \( f(z) \) be given by (71) and the contour \( \gamma_0 \) consists of the union of oriented simple arcs \( \gamma_0 = (\bigcup \gamma_{m,j}) \cup (\bigcup \gamma_{c,j}) \) with distinct end points \( \vec{a}_0 \). Assume that \( \gamma_0 \) only crosses the real axis at \( z_0 = \frac{\mu}{2} \) and \( \gamma_0 \) does not cross the interval \([-T,T]\). Then \( f \in \mathbb{L}(\gamma, \vec{a}, \vec{\beta}) \), where \( \vec{\beta} = (\mu, x, t) \), for all \( x > 0, t > 0, \) and \( \mu > 0 \).

**Proof.** Follows from Definition 4.3 with the logarithmic singularity at \( z_0 = \frac{\mu}{2} \), Lemma 5.1 and Lemma 3.2.

Note the jump of \( f(z) \) is caused by the Schwarz reflection (72) on the real axis and it is linear in \( z \) since \( \Im f \) is a linear function on the real axis (as a limit) near \( \frac{\mu}{2} \) with \( \Im f \left( \frac{\mu}{2} \right) = 0 \) [20]. \( \square \)

Now we can apply all results from Section 4 including the Perturbation theorem 4.1.
The main objects are

\[ K(\alpha_j, \alpha, \mu, x, t) = \frac{1}{2\pi i} \begin{vmatrix}
    f_1 & \frac{d\zeta}{R(\zeta)} & \ldots & \frac{\zeta^{N-1}d\zeta}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \frac{d\zeta}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)}
    \vdots & \vdots & \ddots & \vdots & \ddots & \vdots & \vdots & \ddots & \vdots
    \frac{f_1}{R(\zeta)} & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)}
    \frac{f_{g_1}}{R(\zeta)} & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)}
\end{vmatrix} \]

and

\[ \frac{\partial K(\alpha_j)}{\partial \mu}(\alpha, \mu, x, t) = \frac{1}{2\pi i} \begin{vmatrix}
    f_1 & \frac{d\zeta}{R(\zeta)} & \ldots & \frac{\zeta^{N-1}d\zeta}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \frac{d\zeta}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)}
    \vdots & \vdots & \ddots & \vdots & \ddots & \vdots & \vdots & \ddots & \vdots
    \frac{f_1}{R(\zeta)} & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)} & \ldots & \frac{f_{m,1}}{R(\zeta)}
    \frac{f_{g_1}}{R(\zeta)} & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)} & \ldots & \frac{f_{g_2}}{R(\zeta)}
\end{vmatrix} \]

where \( f_\mu \) is given by (74).

**Theorem 5.3.** \((\mu, x, t)-\text{Perturbation theorem})

Consider a simple contour \( \gamma_0 \) consisting of a finite union of oriented simple arcs \( \gamma_0 = \bigcup \gamma_{m,j} \cup \bigcup \gamma_{c,j} \) with the distinct end points \( \alpha_0 \) and depending on parameters \( \beta_0 = (\mu_0, x_0, t_0) \) (see Figure 1). Assume \( \alpha_0 \) and \((\mu_0, x_0, t_0)\) satisfy a system of equations

\[ K(\alpha_0, (\mu_0, x_0, t_0)) = 0, \]

and \( f \) is given by (71). Let \( \gamma = \gamma(\alpha, \mu) \) be the contour of a RH problem which seeks a function \( h(z) \) which satisfies the following conditions

\[
\begin{align*}
    h_+(z) + h_-(z) &= 2W_j, \quad \text{on } \gamma_{m,j}, \quad j = 0, 1, \ldots, N, \\
    h_+(z) - h_-(z) &= 2\Omega_j, \quad \text{on } \gamma_{c,j}, \quad j = 1, \ldots, N, \\
    h(z) + f(z) &\text{ is analytic in } \mathbb{C} \setminus \gamma,
\end{align*}
\]

where \( \Omega_j = \Omega_j(\alpha, \mu, x, t) \) and \( W_j = W_j(\alpha, \mu, x, t) \) are real constants whose numerical values will be determined from the RH conditions. Assume that there is a function \( h(z, \alpha_0, \mu_0, x_0, t_0) \) which satisfies (58) and suppose \( \frac{h'(z, \alpha_0, \mu_0, x_0, t_0)}{R(z, \alpha_0)} \neq 0 \) for all \( z \) on \( \gamma \).

Then the solution \( \alpha = \alpha(\mu, x, t) \) of the system

\[ K(\alpha, (\mu, x, t)) = 0 \]
and \( h(z, \vec{a}(\mu, x, t), \mu, x, t) \) which solves \((58)\) are uniquely defined and smooth in \((\mu, x, t)\) in some neighborhood of \((\mu_0, x_0, t_0)\).

Moreover, \( \Omega_j(\mu, x, t) = \Omega_j(\vec{a}(\mu, x, t), \mu, x, t) \), and \( W_j(\mu, x, t) = W_j(\vec{a}(\mu, x, t), \mu, x, t) \) are defined and smooth in \((\mu, x, t)\) in some neighborhood of \((\mu_0, x_0, t_0)\).

Furthermore, for \( k \geq 1 \):

\[
\frac{\partial \alpha_j}{\partial \mu} (\mu, x, t) = -\frac{2\pi i \partial K(\alpha_j, \vec{a}(\mu, x, t))}{D(\vec{a}(\mu, x, t), \mu, x, t) f(\gamma(\mu)) \frac{\partial f}{\partial \mu}(\xi - z) R(\xi, \vec{a})} d\xi,
\]

\[
\frac{\partial h}{\partial \mu} (z, \vec{a}, \mu, x, t) = \frac{R(z, \vec{a})}{2\pi i} \int_\gamma(\mu) \frac{\partial f}{\partial \mu}(\xi, \mu, x, t) d\xi,
\]

where \( z \) is inside of \( \gamma \),

\[
\frac{\partial \Omega_j}{\partial \mu} (\mu, x, t) = \frac{1}{D} \begin{vmatrix}
\frac{d \gamma_1}{d \zeta} & \cdots & \frac{\zeta^{N-1} d \zeta}{d \zeta} & \cdots & \frac{d \gamma_1}{d \zeta} \\
\vdots & \ddots & \vdots & \ddots & \vdots \\
\frac{d \gamma_N}{d \zeta} & \cdots & \frac{\zeta^{N-1} d \zeta}{d \zeta} & \cdots & \frac{d \gamma_N}{d \zeta} \\
\end{vmatrix},
\]

\[
\frac{\partial W_j}{\partial \mu} (\mu, x, t) = \frac{1}{D} \begin{vmatrix}
\frac{d \gamma_1}{d \zeta} & \cdots & \frac{\zeta^{N-1} d \zeta}{d \zeta} & \cdots & \frac{d \gamma_1}{d \zeta} \\
\vdots & \ddots & \vdots & \ddots & \vdots \\
\frac{d \gamma_N}{d \zeta} & \cdots & \frac{\zeta^{N-1} d \zeta}{d \zeta} & \cdots & \frac{d \gamma_N}{d \zeta} \\
\end{vmatrix},
\]

where \( R(\xi) = R(\xi, \vec{a}) \).
Figure 6: The jump contour in the case of genus 2 with complex conjugate symmetry in the notation of [20].

Remark 5.4. In the genus 2 case \((N = 1)\), taking into account the symmetry

\[
\alpha_{2j} = \overline{\alpha}_{2j+1}, \quad j = 0, 1, 2,
\]

in the notation of [20] where the numeration of \(\alpha\)'s was different, \(\hat{\gamma}_m\) is a loop around the main arc \([\alpha_2, \alpha_4]\) and its complex conjugate \([\alpha_3, \alpha_5]\), and \(\hat{\gamma}_c\) is loop around the complementary arc \([\alpha_0, \alpha_2]\) and its complex conjugate \([\alpha_1, \alpha_3]\) (see Fig. 6).

Then the equations (89-93) read

\[
\begin{vmatrix}
\mathbf{f}_{\hat{\gamma}_m} \frac{\partial \Omega}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_m} \frac{\partial f_{\hat{\gamma}_c}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_m} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}_c} \frac{\partial \Omega}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_c} \frac{\partial f_{\hat{\gamma}_c}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_c} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}(\mu)} & \mathbf{f}_{\hat{\gamma}(\mu)} & \mathbf{f}_{\hat{\gamma}(\mu)} \\
\end{vmatrix} = 0,
\]

(95)

\[
\frac{\partial \alpha_j}{\partial \mu}(\mu, x, t) = -\frac{2\pi i}{D} \begin{vmatrix}
\mathbf{f}_{\hat{\gamma}_m} \frac{\partial \Omega}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_m} \frac{\partial f_{\hat{\gamma}_c}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_m} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}_c} \frac{\partial \Omega}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_c} \frac{\partial f_{\hat{\gamma}_c}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_c} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}(\mu)} & \mathbf{f}_{\hat{\gamma}(\mu)} & \mathbf{f}_{\hat{\gamma}(\mu)} \\
\end{vmatrix},
\]

(96)

\[
\frac{\partial h}{\partial \mu}(z, \overline{\alpha}, \mu, x, t) = \frac{R(z, \overline{\alpha})}{2\pi i} \int_{\hat{\gamma}(\mu)} \frac{\partial f_{\hat{\gamma}_m}(\xi, \mu, x, t)}{(\xi - z) R(\xi, \overline{\alpha})} d\xi,
\]

(97)

where \(z\) is inside of \(\hat{\gamma}\),

\[
\frac{\partial \Omega}{\partial \mu}(\mu, x, t) = -\frac{1}{D} \begin{vmatrix}
\mathbf{f}_{\hat{\gamma}_m} \frac{\partial f_{\hat{\gamma}_m}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_m} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}_c} \frac{\partial f_{\hat{\gamma}_c}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_c} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}(\mu)} & \mathbf{f}_{\hat{\gamma}(\mu)} \\
\end{vmatrix},
\]

(98)

\[
\frac{\partial W}{\partial \mu}(\mu, x, t) = -\frac{1}{D} \begin{vmatrix}
\mathbf{f}_{\hat{\gamma}_m} \frac{\partial f_{\hat{\gamma}_m}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_m} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}_c} \frac{\partial f_{\hat{\gamma}_c}}{\partial \xi} & \mathbf{f}_{\hat{\gamma}_c} \frac{\partial \xi}{\partial \xi} \\
\mathbf{f}_{\hat{\gamma}(\mu)} & \mathbf{f}_{\hat{\gamma}(\mu)} \\
\end{vmatrix},
\]

(99)
where \( \alpha_j = \alpha_j(\mu, x, t) \), \( R(\zeta) = R(\zeta, \vec{\alpha}(\mu, x, t)) \), \( f(\zeta) = f(\zeta, \mu, x, t) \), \( \frac{\partial f}{\partial \mu}(\zeta) = \frac{\partial f}{\partial \mu}(\zeta, \mu) \) and

\[
D = D(\mu, x, t) = \left| \frac{f_{\zeta m} \frac{d \zeta}{d \zeta}}{f_{\zeta c} \frac{d \zeta}{d \zeta}} \right| .
\] (100)

Remark 5.5. The perturbation theorem 5.3 guarantees that the solution of the RHP (88) is uniquely continued with respect to external parameters. Additional sign conditions on \( \Im h \) need to be satisfied, for \( h \) to correspond to an asymptotic solution of NLS as in [20]. The sign conditions have to be satisfied near \( \gamma \) and additionally on a semiinfinite complementary arcs connecting the end points of \( \gamma \) to \( \infty \).

6 Appendix

6.1 Alternative formulation of the scalar RHP: function \( g' \)

An alternative approach is to start with a scalar RHP on \( g' \) which we formulate as the following: given an ordered sequence of distinct points \( \vec{\alpha} = (\alpha_0, ..., \alpha_{2N+1}) \) and oriented arcs \( \gamma_{m,j}, \gamma_{c,j} \) connecting them, find a function \( g'(z) \) so the following Riemann-Hilbert conditions are satisfied

\[
\begin{cases}
g'_+(z) + g'_-(z) = f'(z), \text{ on } \gamma_{m,j}, \quad j = 0, 1, ..., N, \\
g'(z) \text{ is analytic in } \mathbb{C} \setminus \bigcup \gamma_{m,j}, \\
g'(z) \sim O(z^{-2}), \text{ as } z \to \infty,
\end{cases}
\] (101)

where by \( g'_\pm \) are denoted the limiting values of \( g' \) approaching the contour from the positive and the negative sides respectively.

If \( g(z) \) satisfies (101) then \( g(z) \) satisfies the following Riemann-Hilbert conditions

\[
\begin{cases}
g_+(z) + g_-(z) = f(z) + W_j, \text{ on } \gamma_{m,j}, \quad j = 0, 1, ..., N, \\
g_+(z) - g_-(z) = \Omega_j, \text{ on } \gamma_{c,j}, \quad j = 0, 1, ..., N, \\
g(z) \text{ is analytic in } \mathbb{C} \setminus \gamma,
\end{cases}
\] (102)

where \( \gamma = (\bigcup \gamma_{m,j}) \cup (\bigcup \gamma_{c,j}) \) is an oriented simple curve called the jump contour. The constants \( W_j, \Omega_j \) are found from the condition of \( g(z) \) being analytic at \( \infty \).

The conditions (101) arguably are more natural since they do not involve unknown constants \( W_j \) and \( \Omega_j \) which are found from the behavior at infinity.

6.2 Evaluation of a simplified integral for Lemma 3.4

Let \( \beta_1 = \mu \), and all other parameters \( \beta_2, \ldots \) are constants and

\[
f(z, \mu) = c(\mu)(z - z_0) \log(z - z_0),
\]

where \( z_0 = z_0(\mu) \), \( c(\mu) \) are continuously differentiable in \( \mu \) and some branch of the logarithm is chosen with the branchcut passing from \( z_0 \) to infinity and not intersecting the contour \([z_1, z_0(\mu)) \cup (z_0(\mu), z_2]\).
Consider the contour integral through $z_0$ with some fixed $z_1$ and $z_2$

$$I_1(\mu) = \int_{[z_1, z_0(\mu)] \cup [z_0, z_2]} f(\xi, \mu) d\xi$$

$$= c(\mu) \int_{[z_1, z_0] \cup [z_0, z_2]} (\xi - z_0) \log(\xi - z_0) d\xi$$

and after the change of variables $y = \xi - z_0$

$$= c(\mu) \int_{[z_1 - z_0, 0] \cup [0, z_1 - z_0]} y \log y dy.$$

Thus

$$I_1(\mu) = c(\mu) \left[ \left( \log(z_2 - z_0) - \frac{1}{2} \right) \frac{(z_2 - z_0)^2}{2} - \left( \log(z_1 - z_0) - \frac{1}{2} \right) \frac{(z_1 - z_0)^2}{2} \right].$$

The derivative of $I_1(\mu)$ in $\mu$ is

$$\frac{dI_1}{d\mu}(\mu) = \frac{c'(\mu)}{c(\mu)} I_1 + c(\mu) \left[ \left( \log(z_2 - z_0) - \frac{1}{2} \right) (z_2 - z_0)(-z_0') - \left( \log(z_1 - z_0) - \frac{1}{2} \right) (z_1 - z_0)(-z_0') \right].$$

On the other hand

$$\frac{\partial f(z, \mu)}{\partial \mu} = c'(\mu) \frac{f(z, \mu)}{c(\mu)} + c(\mu) \left[ \log(z - z_0)(-z_0') + (-z_0') \right]$$

and

$$\int_{[z_1, z_0] \cup [z_0, z_2]} \frac{\partial f(\xi, \mu)}{\partial \mu} d\xi = \frac{c'(\mu)}{c(\mu)} I_1(\mu) - (z_2 - z_1)c(\mu) z_0'(\mu) - c(\mu) z_0' \int_{[z_1 - z_0, 0] \cup [0, z_2 - z_0]} \log y dy,$$

where $y = \xi - z_0$

$$= \frac{c'(\mu)}{c(\mu)} I_1(\mu) - c(\mu) z_0' \left[ (z_2 - z_0) \log(z_2 - z_0) - (z_1 - z_0) \log(z_1 - z_0) \right] = \frac{\partial I_1}{\partial \mu}.$$ 

So

$$\frac{dI_1}{d\mu}(\mu) = \int_{[z_1, z_0] \cup [z_0, z_2]} \frac{\partial f(\xi, \mu)}{\partial \mu} d\xi$$

the derivative and the integral can be interchanged.

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