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Separating the shortwave and longwave components of greenhouse gas radiative forcing

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Abstract
Many important greenhouse gases (including water vapour, carbon dioxide, methane and ozone) absorb solar radiation. When gas concentrations change, this absorption exerts a radiative forcing that modifies the thermal infrared (‘longwave’) radiative forcing which is predominant for most gases (ozone being a major exception). The nature of the solar forcing differs from the longwave forcing in several ways. For example, the sign of the instantaneous solar forcing can differ between the tropopause and top-of-atmosphere, and the sign can differ between gases. In addition, a significant part of the solar forcing can be manifested in the longwave, following stratospheric temperature adjustment, which can counteract or enhance the instantaneous solar forcing. Here the nature of solar forcing is examined via a mixture of idealised and more realistic calculations, which consider the effect of perturbations in carbon dioxide, methane and ozone. An apparent contradiction in the sign of the solar forcing of carbon dioxide is resolved; it is shown to be negative, reducing the net carbon dioxide forcing by about 2.3%. The relevance of this work to the effective radiative forcing concept is also discussed.

KEYWORDS
carbon dioxide, methane, ozone, radiative forcing, solar radiation

1 | INTRODUCTION

As well as absorbing thermal infrared radiation (‘longwave [LW]’, hereafter), greenhouse gases (hereafter GHGs, including H$_2$O, CO$_2$, O$_3$ and CH$_4$) absorb incoming solar radiation at near-infrared and sometimes visible and ultra-violet wavelengths (e.g., Gordon et al., 2017). We refer to this as shortwave (SW), often taken to be wavelengths less than 4 μm, acknowledging that some solar radiation is incident at longer wavelengths.

Many papers consider GHG LW radiative forcing in detail (e.g., Ramaswamy et al., 2018) and several explore mechanisms by which this forcing is manifested; Dufresne et al. (2020) presented an elegant demonstration of the relative contributions of increased atmospheric opacity, and the related change in emission height to CO$_2$ forcing. However, while GHG shortwave absorption, especially by H$_2$O, CO$_2$ and O$_3$, has long been included in climate model calculations (e.g., Manabe & Wetherald, 1967), its radiative forcing role has been relatively neglected, with the exception of O$_3$. In addition, the radiative forcing literature indicates apparent contradictions in the sign of the SW forcing. An intercomparison of climate model radiative transfer codes (Collins et al., 2006) demonstrated that, at that time,
many ignored GHG SW absorption beyond those mentioned above; Pincus et al. (2020) report progress in recent years.

This letter aims to resolve apparent contradictions in the nature of SW GHG radiative forcing. After reviewing current understanding, idealised and more realistic model calculations illustrate the SW processes in the framework of the instantaneous radiative forcing (IRF) and radiative forcing including stratospheric temperature adjustment (henceforth RF) (e.g., Myhre et al., 2013). The relevance to Effective Radiative Forcing (ERF) (e.g., Myhre et al., 2013) is also discussed.

2 | CURRENT UNDERSTANDING

Hansen et al. (1981) presented calculations of top-of-atmosphere (TOA) LW and SW forcing due to a CO2 doubling from 300 ppm. IRFSW was 0.1 W m\(^{-2}\), compared to IRFLW of 2.4 W m\(^{-2}\). The TOA RF\(_{\text{SW}}\) (i.e., forcing including stratospheric temperature adjustment) of 0.1 W m\(^{-2}\) is unchanged from its IRF value; RF\(_{\text{LW}}\) is 3.8 W m\(^{-2}\), significantly higher than RF\(_{\text{LW}}\). Hence, Hansen et al. (1981) found a positive SW forcing of 2.6% or 4% of the net (LW + SW) forcing, depending on whether RF or IRF is considered. The view that CO\(_2\) IRFSW is positive re-emerged in the ERF framework, which takes a TOA perspective (e.g., fig. 14-6 of Ramaswamy et al., 2018).

By contrast, Cess et al. (1993) reported that CO\(_2\) IRFSW was negative, and about 6% of IRF\(_{\text{LW}}\) (for a doubling from 330 ppm). This view became established, although not all studies have found a negative CO\(_2\) tropopause IRFSW (Forster et al., 2001). The apparent contradiction is because Cess et al. (1993) defined forcing at the tropopause; Hansen et al. (1981) chose TOA. This still leaves a question as to which perspective is of most value, and whether they can be reconciled.

Myhre et al. (1998) also found a negative IRFSW for CO\(_2\) of order 4% (−0.11 W m\(^{-2}\)) of RF\(_{\text{LW}}\), for a doubling from 278 ppm. However, and of importance here, the additional SW absorption warms the stratosphere (relative to the LW-only case). In the RF framework, Myhre et al. (1998) calculated that this warming led to a positive tropopause RF\(_{\text{LW}}\) (0.05 W m\(^{-2}\)); thus, the net RF due to SW forcing (−0.06 W m\(^{-2}\)) is about half IRFSW.

For increased concentrations of stratospheric H\(_2\)O, Forster and Shine (2002) (see also Forster et al., 2001; Myhre et al., 2007, 2009) found a negative tropopause IRFSW, offsetting about 20% of RF\(_{\text{LW}}\).

Etminan et al. (2016) presented IRFSW calculations for methane; the tropopause IRFSW (for a 750–1800 ppb perturbation) was positive and 6% of the total forcing; accounting for the effect of warming of the stratosphere due to the additional SW absorption on RF\(_{\text{LW}}\), methane’s SW forcing enhanced the LW-only RF by 15%.

Ozone is a distinct GHG because of its strong absorption of ultraviolet (and, to some extent, visible) radiation. Its SW forcing has long been recognised (e.g., Ramaswamy et al., 1992; Ramaswamy and Bowen, 1994; Hauglustaine et al., 1994) but, to our knowledge, the specific role of SW absorption in modifying RF\(_{\text{LW}}\) has not been isolated. The differing TOA and tropopause perspectives have been indicated, but not fully explained, by Michou et al. (2020); they found the signs of TOA ERFSW and ERFLW for stratospheric ozone depletion were opposite to the tropopause RF\(_{\text{LW}}\) and RF\(_{\text{SW}}\) calculations of Checa-Garcia et al. (2018).

Taken together, these studies show that the apparent SW forcing depends on whether a tropopause or TOA perspective is taken, and whether its impact on stratospheric temperature adjustment (and hence on RF\(_{\text{LW}}\)) is accounted for. They also show that the sign of SW forcing varies among gases, sometimes enhancing and sometimes opposing RF\(_{\text{LW}}\). This letter constructs a framework to better understand SW forcing and to resolve apparent contradictions. This stresses that judging the importance of SW forcing via IRFSW, RF\(_{\text{SW}}\) or ERFSW alone gives a misleading impression.

The impact of SW forcing on the LW via stratospheric temperatures (in the RF framework) and rapid (tropospheric and stratospheric) temperature adjustments (in the ERF framework) must be accounted for to give a correct impression of the size and sometimes the sign of SW forcing. In the ERF framework, other adjustments (e.g., in clouds and water vapour) driven by SW processes can also impact on ERFSW; similarly, changes driven by LW processes can impact on ERFSW (e.g., Donohoe et al., 2014).

TOA IRFSW must be positive (consistent with Hansen et al. (1981)) for GHG concentration increases. Additional shortwave absorption always decreases planetary albedo. At the tropopause, the situation is unclear, unless the gas perturbation is solely in the stratosphere. Stratospheric absorption deprives the troposphere of radiation giving a negative IRFSW. However, additional tropospheric absorption decreases the albedo of the troposphere-surface system giving a positive IRFSW. Hence the tropopause IRFSW can have either sign.

Etminan et al. (2016) show that the sign of CO\(_2\) and CH\(_4\) tropopause IRFSW varies with wavelength. They also show that the sign depends on the intensity of absorption features (for strong absorption features, additional absorption is mostly in the stratosphere, leading to a negative IRFSW), and overlap with strong near-IR water vapour bands; if these bands are saturated in the troposphere, additional absorbers exert little influence on the upwelling irradiance at the tropopause. By contrast,
additional absorption in relatively transparent windows between these bands can lead to a positive IRF\textsubscript{SW}. This is especially so for cloudy skies, as tropospheric albedo is higher, and the importance of absorption of upwelling radiation by the gas is enhanced; Etminan et al. (2016) find the wavelength-integrated CH\textsubscript{4} IRF\textsubscript{SW} at the tropopause is negative for clear skies but positive when clouds are included (although Collins et al., 2018 found it to be positive in both cases). Etminan et al. (2016) found it was negative in both cases for CO\textsubscript{2}.

As noted above, IRF alone cannot constrain the effect of shortwave absorption; within the RF framework, the impact of increased stratospheric SW absorption on LW forcing must be considered. Section 3 uses idealised calculations to illustrate this.

## 3 | IDEALISED CALCULATIONS WITH NO INSTANTANEOUS LONGWAVE RADIATIVE FORCING

The idealised calculations of SW forcing employ the LW code of Shine and Myhre (2020) and the SW code of Slingo and Schrecker (1982) with updated gaseous absorption coefficients from Chagas et al. (2001). Calculations use a global-mean atmospheric profile (temperature, humidity, ozone and clouds) from Freckleton et al. (1998) with a global-mean tropopause of 128.6 hPa, global-mean insolation (solar zenith angle of 60° for 12 h) and a spectrally-constant surface albedo of 0.06 (representing a sea surface). The forcing due to a grey absorber which is added only to near-IR bands (wavelengths greater than 1 μm) is computed; that is, there is no IRF\textsubscript{LW}. The grey absorber has an absorption coefficient of $4 \times 10^{-4}$ m\textsuperscript{2} kg\textsuperscript{-1} and a constant mass mixing ratio of 0.0005 kg kg\textsuperscript{-1}, giving an optical depth of 0.003 when the absorber is in the stratosphere only and 0.024 when at all altitudes. RF is computed by a standard time-stepping procedure that adjusts stratospheric temperature until the LW + SW heating rates return to global-mean radiative equilibrium.

In Idealised Example 1 (Table 1), the grey absorber is in the stratosphere only. As expected from Section 2, IRF\textsubscript{SW} is positive at TOA, and negative at the tropopause.

This causes additional stratospheric absorption of solar radiation; in this case, the TOA forcing is +0.29 W m\textsuperscript{-2}, the tropopause forcing is −0.54 W m\textsuperscript{-2}, giving a convergence of 0.83 W m\textsuperscript{-2}. The consequent warming of the stratosphere increases LW emission to space and the troposphere. Increased upward TOA irradiance constitutes a negative LW forcing; increased downward tropopause irradiance constitutes a positive LW forcing. Thus, RF\textsubscript{LW} is opposite in sign and comparable in size to IRF\textsubscript{SW} at both TOA and tropopause (Table 1).

It is initially surprising that at TOA, RF\textsubscript{LW} (−0.4 W m\textsuperscript{-2}) is larger in magnitude than IRF\textsubscript{SW}. However, increased LW emission from the stratosphere due to the adjustment (0.4 W m\textsuperscript{-2} upwards at TOA and 0.43 W m\textsuperscript{-2} downwards at the tropopause) is consistent, as it should be, with the 0.83 W m\textsuperscript{-2} convergence of SW radiation. In this case, RF\textsubscript{LW} is nearly equal at TOA and tropopause; the SW effect on RF\textsubscript{LW} could be estimated by partitioning the convergence of SW radiation in this way. This is because the grey absorber is at all stratospheric levels. When placed in the topmost layer only (at 1 hPa), TOA RF\textsubscript{LW} is about 6 times larger than the tropopause RF\textsubscript{LW}. When placed only in the layer closest to the tropopause, the tropopause RF\textsubscript{LW} is about double TOA RF\textsubscript{LW}.

The net forcing, RF\textsubscript{NET}, at TOA and tropopause is now equal (−0.11 W m\textsuperscript{-2}), as is required following stratospheric temperature adjustment. In this example, because SW absorption deprives the surface-troposphere system of energy, RF\textsubscript{NET} is negative but, because of the compensatory effect of increased stratospheric LW emission, it is only 20% of the value inferred from the tropopause IRF\textsubscript{SW}.

This illustrates how IRF\textsubscript{SW} differs in sign between the TOA and tropopause perspectives (i.e., there is no contradiction in the literature) and also illustrates how SW absorption cannot be judged from IRF\textsubscript{SW} alone; the effect on RF\textsubscript{LW} must be considered. Once RF\textsubscript{LW} is included, there is no ambiguity in the sign of RF\textsubscript{NET} and it agrees at TOA and tropopause.

In Idealised Example 2 (Table 2) the grey absorber is present at all altitudes. IRF\textsubscript{SW} is now positive at both TOA and tropopause, because of increased tropospheric absorption of solar radiation. Because stratospheric convergence of SW radiation is only slightly affected by

| W m\textsuperscript{-2} | Instantaneous | Adjusted |
|-------------------------|--------------|----------|
|                         | LW  | SW   | Net | LW  | SW   | Net |
| TOA                     | 0.0 | +0.29| +0.29| −0.40| +0.29| −0.11|
| Tropopause              | 0.0 | −0.54| −0.54| +0.43| −0.54| −0.11|


tropospheric absorption (i.e., 1.50–0.67 = 0.83 W·m⁻² as in Example 1), RF_LW from stratospheric temperature adjustment is almost identical to Table 1. In this case, RF_LW is a smaller proportion of IRF_SW, and RF_NET is positive.

Again, this example illustrates that forcing due short-wave absorption cannot be judged by IRF_SW alone, although in this case the sign of IRF_SW is the same at TOA and tropopause and consistent with RF_NET.

### 4 | MORE REALISTIC CALCULATIONS FOR CARBON DIOXIDE, METHANE AND OZONE

The role of SW forcing in more realistic cases is calculated using the more sophisticated configuration of Checa-Garcia et al. (2018). RF is calculated on a 5° × 5° horizontal grid; stratospheric temperature adjustment is calculated using the fixed-dynamical heating method. It uses the SOCRATES radiative transfer code (Walters et al., 2019), using the Met Office Earth System Model configuration: 9 LW bands (wavenumbers 1–2995 cm⁻¹) and 6 SW bands (wavenumbers 1–50,000 cm⁻¹). Unlike Section 3, IRF has both LW and SW components. We perform calculations with both LW and SW components, and then repeat them with only the LW component active (‘LW only’ in the tables). The difference between these yields the total RF_SW forcing, including its impact on RF_LW via stratospheric temperature adjustment.

SOCRATES is regularly updated to reflect its performance in radiation code intercomparisons (e.g., Pincus et al., 2015) and updated spectral data. Walters et al. (2019) (their sec. 3.2.1) document significant improvements in the version used here, relative to a high-spectral resolution code which was compared with other benchmark codes in Pincus et al. (2020).

The example GHGs (CO₂, CH₄, O₃) are the ones most widely discussed in earlier work (Section 2); their different behaviours should guide how other GHGs would behave. CO₂ has intense SW stratospheric absorption so that its tropopause IRF_SW is negative; methane is weaker giving a positive tropopause IRF_SW; ozone is unusual as RF_SW and RF_LW are comparable. Etminan et al. (2016) demonstrate that nitrous oxide’s RF_SW is much smaller than gases considered here. The results presented here are highly relevant to the ERF framework. Stratospheric temperature adjustment is the largest adjustment in ERF calculations for CO₂ (Smith et al., 2018) and ozone (Skeie et al., 2020); for methane, adjustments are small when RF_SW is neglected, but are more important when it is included (Etminan et al., 2016).

Calculations use multi-year (2000–2009) monthly-mean averages of temperature, water vapour, clouds and surface albedo from ERA-Interim (Checa-Garcia et al., 2018). The source of ozone fields is described below. Calculations include the impact of tropospheric scattering of solar radiation by clouds and the surface on absorption of SW radiation in the stratosphere; this has been shown (Section 2) to be important in quantifying methane’s IRF_SW (Collins et al., 2018; Etminan et al., 2016). To demonstrate the role of SW forcing, methane and CO₂ calculations are presented for January, as there is a relatively small seasonal dependence; for ozone, where seasonal variations larger, results are presented as annual-means derived from monthly-mean calculations.

Table 3 (and Figure 1) shows the forcing for CO₂ increasing from 278 to 417 ppm; the IRF_SW is +5.6% of IRF_NET at TOA and –5.2% at the tropopause, consistent with earlier literature (Section 2). The effect of stratospheric temperature adjustment on RF_NET is much larger at TOA (increasing it by 70%) than at the tropopause (decreasing it by 8%), consistent with earlier literature. This means that RF_SW (which is unchanged from IRF_SW because of the weak impact of stratospheric temperature change) is a smaller component of RF_NET at TOA (3.3%) and a slightly larger tropopause component (–5.6%). As required, RF_NET now agrees at TOA and tropopause but, as shown in Table 3 and Figure 1 (left), the apparent SW attribution differs in sign between these levels.

However, this does not account for the role of SW absorption in temperature adjustment. This can be assessed by calculating RF_NET due to LW processes alone (Etminan et al., 2016; Myhre et al., 1998). The lower two rows in Table 3 shows this LW-only RF_NET is 0.05 W·m⁻² greater than the full RF_NET. This 0.05 W·m⁻² reduction in RF_NET robustly indicates the

| W m⁻² | Instantaneous | | Adjusted |
|-------|--------------|---|---|
|       | LW | SW | Net | LW | SW | Net |
| TOA   | 0.0 | +1.50 | +1.50 | –0.39 | +1.50 | +1.11 |
| Tropopause | 0.0 | +0.67 | +0.67 | +0.43 | +0.67 | +1.11 |
TABLE 3  Global-mean instantaneous and adjusted radiative forcing (W m\(^{-2}\)) for January at the top-of-atmosphere and tropopause for an increase in CO\(_2\) from 278 to 417 ppm

| W m\(^{-2}\)      | Instantaneous | Adjusted |
|-------------------|--------------|----------|
|                   | LW | SW | Net | LW | SW | Net |
| TOA               | +1.18 | +0.07 | +1.25 | +2.07 | +0.07 | +2.14 |
| Tropopause        | +2.44 | -0.12 | +2.32 | +2.26 | -0.12 | +2.14 |
| TOA LW only (SW impact) | +1.18 | 0.0 | +1.18 | +2.19 (−0.12) | 0.0 (+0.07) | +2.19 (−0.05) |
| Tropopause LW only (SW impact) | +2.44 | 0.0 | +2.44 | +2.19 (+0.07) | 0.0 (−0.12) | +2.19 (−0.05) |

Note: The third and fourth rows show the case for LW-only calculations. The values in parentheses are the component of the full LW + SW forcing that is attributed to SW forcing; these are derived by differencing row 1 from row 3 for TOA and row 2 from row 4 for the tropopause.

FIGURE 1 (left) radiative forcing (in W m\(^{-2}\)) at the top-of-atmosphere (TOA) and tropopause (TROP) separated into longwave, shortwave and net forcing, for the CO\(_2\) and CH\(_4\), and stratospheric and whole atmosphere O\(_3\) perturbations described in the text. (right) radiative forcing separated into longwave, shortwave and net forcing when the component of LW forcing due to temperature adjustment resulting from the SW forcing is attributed to the SW forcing. In this case, TOA and TROP forcings are identical.

TABLE 4  Global-mean instantaneous and adjusted radiative forcing (W m\(^{-2}\)) for January at the top-of-atmosphere and tropopause for an increase in CH\(_4\) from 725 to 1450 ppb

| W m\(^{-2}\)      | Instantaneous | Adjusted |
|-------------------|--------------|----------|
|                   | LW | SW | Net | LW | SW | Net |
| TOA               | +0.37 | +0.07 | +0.44 | +0.33 | +0.07 | +0.40 |
| Tropopause        | +0.36 | +0.01 | +0.37 | +0.39 | +0.01 | +0.40 |
| TOA LW only (SW impact) | +0.37 | 0.0 | +0.37 | +0.37 (−0.04) | 0.0 (+0.07) | +0.37 (+0.03) |
| Tropopause LW only (SW impact) | +0.37 | 0.0 | +0.37 | +0.37 (+0.02) | 0.0 (+0.01) | +0.37 (+0.03) |

Note: The third and fourth rows show the case for LW-only calculations. The values in parentheses are then the component of the full LW + SW forcing that is attributed to SW forcing.

total impact of SW forcing, accounting for the direct effect via IRF\(_{SW}\), and its indirect impact on RF\(_{LW}\) via stratospheric temperature change. RF\(_{NET}\) decreases by 2.3% at both TOA and tropopause (see Figure 1 [right]). As in Section 2, the contribution of the direct IRF\(_{SW}\), and its impact on RF\(_{LW}\), differs depending on TOA and tropopause perspectives. In this case the positive TOA IRF\(_{SW}\) gives an incorrect perception of the sign of SW absorption. Tropopause IRF\(_{SW}\) significantly overemphasises the size of the (negative) SW forcing, as noted by Myhre et al. (1998).

In the Section 3 idealised calculations, the RF\(_{LW}\) due to SW absorption was approximately equal and opposite at TOA and tropopause. For CO\(_2\), the additional SW absorption is mostly in the upper stratosphere; the effect on RF\(_{LW}\) is about 1.7 times higher at TOA than at the
 tropopause (compare the $-0.12$ and $0.07 \text{ W m}^{-2}$ values in parentheses in the adjusted RF$_{\text{LW}}$ column of Table 3).

Table 4 and Figure 1 show results for methane doubling from 725 ppb. We will show elsewhere that the low spectral-resolution version of SOCRATES underestimates methane’s IRF$_{\text{SW}}$; the purpose here is to illustrate processes, rather than to present definitive values for methane RF. In this case, both TOA and tropopause IRF$_{\text{SW}}$ are positive, although TOA IRF$_{\text{SW}}$ is more strongly so. Stratospheric temperature adjustment is small in the LW-only case (Etminan et al., 2016); both IRF$_{\text{NET}}$ and RF$_{\text{NET}}$ are $0.37 \text{ W m}^{-2}$. When SW is included, convergence of SW radiation in the stratosphere ($0.06 \text{ W m}^{-2}$) drives a larger adjustment; IRF$_{\text{NET}}$ and RF$_{\text{NET}}$ differ by about $0.03 \text{ W m}^{-2}$ at TOA and tropopause. Unlike the CO$_2$ case this is not sufficiently strong to reverse the sign of TOA RF$_{\text{SW}}$ (it decreases from $+0.07$ to $+0.03 \text{ W m}^{-2}$) but it significantly enhances tropopause RF$_{\text{NET}}$ (from $0.01$ to $0.03 \text{ W m}^{-2}$) compared to the IRF$_{\text{SW}}$, consistent with Etminan et al. (2016).

Figure 1 shows results for ozone perturbations, taking the CMIP6 case from Checa-Garcia et al. (2018) for stratospheric ozone change (Table 5) and stratospheric and tropospheric ozone change (Table 6) derived from multimodel averages. Forcing is calculated using decadal-mean ozone fields for 2000–2009, relative to 1850–1859.

Ozone differs from CO$_2$ and CH$_4$ because of the perturbation’s more complex morphology, and because SW forcing plays a larger relative role (Figure 1). For stratospheric ozone depletion (Table 5) IRF$_{\text{SW}}$ is negative at TOA (the decreased stratospheric absorption means more SW radiation is reflected), and positive at the tropopause (more radiation is transmitted through the stratosphere). In both cases IRF$_{\text{LW}}$ is negative. The instantaneous divergence of forcing across the stratosphere ($\approx 0.07 \text{ W m}^{-2}$ due to IRF$_{\text{LW}}$ and $\approx 0.22 \text{ W m}^{-2}$ due to IRF$_{\text{SW}}$) drives strong stratospheric cooling. The reduced emission increases TOA RF$_{\text{LW}}$ relative to IRF$_{\text{LW}}$, changing its sign from $-0.045$ to $+0.088 \text{ W m}^{-2}$, and makes tropopause RF$_{\text{LW}}$ more negative ($-0.02$ to $-0.13 \text{ W m}^{-2}$). This reduces TOA RF$_{\text{NET}}$ and changes the sign of tropopause RF$_{\text{NET}}$. Importantly, even though RF$_{\text{NET}}$ is identical at TOA and tropopause, the LW and SW components are of opposite signs, explaining the apparent discrepancy mentioned by Michou et al. (2020) (see Section 2). The TOA RF is most consistent with the ERF perspective.

By comparing with the LW-only case, Table 5 and Figure 1 show the major effect of SW-induced stratospheric cooling on RF$_{\text{LW}}$. Without SW-induced cooling, RF$_{\text{LW}}$ is $-0.06 \text{ W m}^{-2}$ at TOA and tropopause; with it, they are $+0.09$ and $-0.13 \text{ W m}^{-2}$ respectively.

The case with decreasing stratospheric and increasing tropospheric ozone (Table 6, Figure 1) is more complex than the stratosphere-only case. The biggest difference is the positive IRF$_{\text{LW}}$ at both TOA and tropopause, but IRF$_{\text{SW}}$ is also impacted via reduced SW reflection from the troposphere. The instantaneous divergence of forcing
across the stratosphere (0.23 W m⁻² due to LW and 0.24 W m⁻² due to SW) still drives strong stratospheric cooling. The SW forcing, via its impact on stratospheric temperature, increases TOA RF₁, LW by 85% (from 0.19 to 0.34 W m⁻²) and decreases tropopause RF₁, LW by 45% (from 0.19 to 0.11 W m⁻²). Unlike Table 5, RF₁, LW is positive at tropopause and TOA for both the full and LW-only case, as the forcing from tropospheric ozone increases dominates. As shown in Figure 1 (right), RF₁, LW is the dominant contributor to RF_NET, even when the impact of SW forcing on RF₁, LW is accounted for.

5 | CONCLUSIONS

Via idealised and more realistic calculations, the nature of SW radiative forcing has been investigated. It has been shown that even the sign of the SW forcing can differ between top-of-atmosphere and tropopause perspectives, even though, following stratospheric temperature adjustment, the net top-of-atmosphere and tropopause forcings are identical. This indicates that, on its own, the shortwave forcing is not a consistent indicator of its importance in net forcing.

A more consistent view is achieved by considering the impact of SW forcing on LW forcing via stratospheric temperature adjustment. This separation can be achieved by comparing calculations that include and exclude SW forcing.

In this perspective, not only do the top-of-atmosphere and tropopause perspectives agree in the net forcing, but also the partitioning between SW and LW agrees. In the specific case of increased CO₂, SW processes decrease the net forcing at both the top of the atmosphere and tropopause by 2.3%; this resolves an apparent contradiction in the earlier literature that indicated that the sign of the SW forcing differs between these perspectives. For methane, the instantaneous SW tropopause forcing is smaller than the top-of-atmosphere because it is a residual of negative forcing due to increased stratospheric absorption and positive forcing due to decreased tropospheric reflection; including the effect of this SW absorption on stratospheric temperatures achieves a more nuanced but consistent view. For the ozone, again the top-of-atmosphere and tropopause views of the importance of SW and LW components differ significantly unless the SW influence on LW forcing is accounted for.

The most important conclusion here is that in both radiative forcing and effective radiative forcing frameworks, the role of SW forcing, when it arises from atmospheric absorption (rather than scattering), cannot be assessed by considering changes in SW irradiances alone; indeed, even the implied sign of the SW forcing may be incorrect. We have demonstrated that this is the case for stratospheric temperature adjustment. More detailed calculations with ESMs would be needed to understand how other LW rapid adjustments are affected by SW forcings to achieve a more complete view.

AUTHOR CONTRIBUTIONS

Keith Shine: Conceptualization; formal analysis; investigation; methodology; supervision; visualization; writing – original draft; writing – review and editing. Rachael Byrom: Investigation; methodology; writing – original draft; writing – review and editing. Ramiro Checa-Garcia: Investigation; methodology; software; writing – original draft; writing – review and editing.

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