Pulse phase and precession phase resolved spectroscopy of Her X-1: studying a representative Main-On with RXTE

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ABSTRACT

\textbf{Aims.} We performed a detailed pulse phase resolved spectroscopy of the accreting binary X-ray pulsar Her X-1 in the energy range 3.5–75 keV and have established pulse phase profiles for all spectral parameters. For three parameters, the centroid energy of the cyclotron line, the photon index and the flux of the 6.4 keV iron line, we have studied the variation as a function of 35 d phase.

\textbf{Methods.} We analyzed RXTE observations of the Main-On of November 2002 using data from the PCA and the HEXTE instruments. Four different time intervals of about \( \sim 1 \) d duration were selected to provide a good coverage of a complete Main-On. The intervals are centered at 35 d phases 0.03, 0.10, 0.15, and 0.20, respectively.

\textbf{Results.} All spectral parameters show a strong modulation with pulse phase. While the centroid energy of the cyclotron line follows roughly the shape of the pulse profile, showing higher values close to the peak of the X-ray pulse, both the photon index and the iron line intensity exhibit distinct minima around the peak of the X-ray pulse. With respect to variations of the observed profiles with 35 d phase, we find that there is a clear evolution of the shape of the pulse profiles (flux versus pulse phase), a moderate increase of the maximum cyclotron line energy (found around pulse phase 0.7), but no significant evolution of the shape of the pulse phase profiles of the cyclotron line energy, the spectral power law index or the iron line intensity.

\textbf{Conclusions.} The variation of spectral parameters as a function of the pulse phase provides important information about the system: i. the disappearance of the Fe line flux near the highest continuum flux may be an indication of a hollow cone geometry of the accretion structure; ii. The apparent non-dependence of the cyclotron line energy profiles on 35 d phase provides a new possibility to test the model of free precession of the neutron star, proposed to be responsible for the systematic variations in the pulse profiles.

\textbf{Key words.} binaries: general – stars: neutron – X-rays: binaries – X-rays: individuals: Her X-1 – pulsars: individuals: Her X-1

1. Introduction

The X-ray source Her X-1 was discovered in 1972 by Uhuru (Tananbaum et al. 1972), and classified as an accreting X-ray binary. Her X-1 is one of the brightest and most studied persistent binary X-ray pulsars. The distance to the system is \( \sim 7 \) kpc and the masses of the neutron star and its companion are approximately 1.5 \( M_\odot \) and 2.2 \( M_\odot \), respectively (Reynolds et al. 1997). The X-ray flux shows periodic modulation on several different time-scales: pulsations with 1.24 s due to the spin of the neutron star, eclipses due to the 1.7 d orbital period, and a 35 d super-orbital period due to obscuration of the X-ray emitting region by the precessing accretion disk. This 35 d periodicity shows different states: two On-states with the \( \sim 11 \) d Main-On (at phase 0–0.31) and the \( \sim 5 \) d Short-On (at phase 0.57–0.79), separated by two Off-states (see e.g. Giacconi et al. 1973; Gerend & Boynton 1976; Boynton et al. 1980; Scott & Leahy 1999; Scott et al. 2000; Klochkov et al. 2006). The accretion disk is thought to be tilted (with respect to the orbital plane), warped and counter-precessing with a somewhat variable period of \( \sim 35 \) d. The onset of the flux (often identified with 35 d phase 0.0) is called the Turn-On (TO) and corresponds to the transition from the Off-state to the Main-On. At this time, the outer rim of the disk clears the view to the X-ray emitting region close to the polar caps on the surface of the neutron star, leading to an increase in flux. Viceversa, the decline of the flux towards the end of the Main-On is identified with the inner edge of the accretion disk blocking our view to the X-ray emitting regions. A similar cycle is responsible for the Short-On with a maximum flux of roughly one third that of the Main-On.

Her X-1 was also the first accreting X-ray pulsar for which a cyclotron line in the X-ray spectrum has been discovered (Trümper et al. 1978). This absorption-like feature, now referred to as Cyclotron Resonance Scattering Feature (CRSF), is observed around 40 keV and allows to estimate the neutron star’s magnetic field. Applying the formula \( B_{12} = (1+z) E_{cyc}/11.6 \) keV (where \( B_{12} \) is the magnetic field strength in units of \( 10^{12} \) Gauss, \( z \) is the gravitational redshift and \( E_{cyc} \) is the centroid energy of the cyclotron line), the first direct measurement of the magnetic field of a neutron star was achieved (\( \sim 3 \times 10^{12} \) Gauss for Her X-1). The source shows a positive correlation between the centroid energy of the cyclotron line and the bolometric luminosity (Staubert et al. 2007; Vasco et al. 2011).

Another observational feature of Her X-1 is the evolution of the pulse profile as function of energy and time, or better: 35 d phase. For both the Main-On and the Short-On, a systematic variation in the shape of the pulse profiles is found. These changes consist of the disappearing of some features in the pulse profile with an increase in energy and/or time. These variations repeat systematically on a time scale of \( \sim 35 \) d. These changes have been well known for some time (see e.g. Deeter et al. 1998; Scott et al. 2000; Trümper et al. 1980). Recently, Staubert et al. (2010a,b, 2012) on the basis of RXTE observations, gave a quantitative description of these changes by providing a template, which can be used to predict the shape of the profile in the 9–
13 keV energy range for any 35 d phase during the Main-On. An extension of this work for other energy ranges and including the Short-On is in preparation.

Trümper et al. (1986), on the basis of observations by Exosat, had proposed that free precession of the neutron star is responsible for those pulse profile changes, due to the variation of the angle of the line of sight to the X-ray emitting regions on the surface of the neutron star. For Her X-1, the dependence of the spectral parameters on pulse phase have been observed by several missions. Early examples are Pravdo et al. (1978) with OSO 8, Pravdo et al. (1979) with HEAO-1 and McCray et al. (1982) with HEAO 2, the Einstein Observatory, Voges et al. (1982) with the MPE/ATI balloon experiment, Kahabka (1987) with Exosat, Kunz et al. (1996) with Mir-HEXE. Later examples are Endo et al. (2000) with ASCA, Zane & Ramsay (2001) with XMM-Newton, Lutovinov et al. (2000) with the ART-P telescope on board of GRANADA and Klochkov et al. (2008) with INTEGRAL.

In this work we present a deep investigation of the variation of key spectral parameters as a function of pulse phase with the finest resolution so far and, for the first time, a discussion of a possible evolution with 35 d phase. Changes of the spectral parameters with pulse phase are quite common among accreting X-ray pulsars and are generally attributed to a change in the viewing angle of the accretion region (see e.g., Kreykenbohm et al. 2004 and references therein). Here, we focus on the variation of the centroid energy $E_{\text{c}}$, of the CRSF around 40 keV, the photon index $\Gamma$, and the iron line characteristics. We also address the question whether our data support the free precession model which has been suggested to explain the variation in shape of the pulse profile.

2. Observations and data analysis

To carry out this analysis we used observations of Her X-1 of the Rossi X-ray Timing Explorer (RXTE) performed in November 2004 and references therein). Here, we focus on the variation of the Main-On of cycle 323. The grey zones represent the four intervals used for this analysis. For each interval the centered 35 d phase is given. Details of the time and 35 d phase intervals are listed in Table I.

Fig. 1. PCA light curve of the Main-On of cycle 323. The grey zones represent the four intervals used for this analysis. For each interval the centered 35 d phase is given. Details of the time and 35 d phase intervals are listed in Table I.

Table 1. Summary of the four intervals of the Main-On of 35 d cycle no. 323. The columns are: the number of the interval, the time interval in MJD, the exposure time, the corresponding 35 d phase interval, and the center of the 35 d phase interval.

| Interval | limits MJD | exposure [ks] | limits 35 d phase | center 35 d phase |
|----------|------------|---------------|-------------------|------------------|
| 1        | 52595.060 – 52596.743 | 25 | 0.007 – 0.056 | 0.03 |
| 2        | 52597.959 – 52598.719 | 25 | 0.090 – 0.112 | 0.10 |
| 3        | 52599.603 – 52600.726 | 25 | 0.137 – 0.161 | 0.15 |
| 4        | 52601.316 – 52602.406 | 28 | 0.186 – 0.217 | 0.20 |

For barycentric and binary corrections orbital parameters and the ephemeris of Staubert et al. (2009) were used: $P_{\text{orb}}=1.700167287$ s, $P_{\text{orb}} = +(2.8 \pm 0.2) \times 10^{-12}$ s s$^{-1}$, $a = \sin i = 13.1831$ lt s and $T_{\text{ej}}=52599.486440$ MJD. To align the pulse profiles, a reference time $t_0 = MJD 52594.86990091$ and a spin period $P_{\text{spin}} = 1.237761809$ s were used. As "pulse phase zero" the sharp-edge feature at the trailing edge of the right shoulder of the main pulse which leads into an eclipse like minimum (Staubert et al. 2009b). Note that these pulse profiles are statistically so accurate that the uncertainties are smaller than the line width.

The PCA was used in the energy range 3.5 – 60 keV and HEXTE in the energy range 20 – 75 keV. The data have been analyzed with XSPEC 12.6.1 using the highecut spectral model.

http://heasarc.gsfc.nasa.gov/docs/xanadu/xspec
http://heasarc.gsfc.nasa.gov/xanadu/xspec/manual/XSmodelHighecut.html

Fig. 2. Mean 9-13 keV pulse profiles of the four time intervals at different 35 d phases (see Table I): Interval 1 in black (solid line), Interval 2 in green (dashed line), Interval 3 in red (dash dotted line) and Interval 4 in blue (dash dot dot dotted line). For general 35 d dependence see Staubert et al. (2010). As timing reference for pulse phase zero we use the "sharp-edge" at the trailing edge of the right shoulder of the main pulse which leads into an eclipse like minimum (Staubert et al. 2009b). Note that these pulse profiles are statistically so accurate that the uncertainties are smaller than the line width.
power law continuum with an exponential cut-off. The differential photon flux \( f(E) \) is described by:

\[
f(E) = \begin{cases} 
  E^{-\Gamma} \cdot \exp \left( \frac{E - E_{\text{cut}}}{E_{\text{fold}}} \right), & \text{if } E \leq E_{\text{cut}} \\
  E^{-\Gamma}, & \text{if } E > E_{\text{cut}} 
\end{cases}
\]

(1)

where \( \Gamma \) is the photon index, \( E_{\text{cut}} \) is the cut-off energy and \( E_{\text{fold}} \) is the e-folding energy. In addition, two line features are included in the model: a multiplicative Gaussian optical depth for the Cyclotron Resonant Scattering Feature (CRSF) and an additive Gaussian emission line for the iron fluorescence feature around 6.5 keV. This model is the same one as used in earlier analyses (see e.g., Coburn et al. [2002] and Staubert et al. [2007]). We have consistently found that the highcut spectral model describes the spectrum of Her X-1 best. The discontinuity of this model at the cut-off energy does not constitute a problem and we have verified that no systematic uncertainties are introduced in estimating the spectral parameters. We used the PCA data up to 60 keV, which has become feasible by new response matrices (PCA response v11.7, 2009 May 11). These high energy data can definitively improve the photon statistics at energies around and beyond the cyclotron line feature. The same choice has been used in previous analyses (see e.g. Ferrigno et al. [2011]) and it was confirmed by Rothschild et al. [2011] with RXTE observations of Cen A. The recorr spectral component was also added to the model to normalize the background. After we had verified, that the energy and width of the iron K line is constant (6.45 keV and 0.5 keV, respectively), we fixed those values in order to minimize the number of free parameters. GoodXenon observational modes of the PCA for high resolution analysis were used to extract light curves and spectra. An additional systematic uncertainty of 1% was added in using the PCA data. To analyze the data, version v6.11 of the HEASOFT software and

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**Fig. 3.** Cyclotron line centroid energy profiles as function of the pulse phase for Interval 1 (top left), Interval 2 (top right), Interval 3 (bottom left) and Interval 4 (bottom right).

**Fig. 4.** The cyclotron line centroid energy as function of pulse phase for the four different 35 d phase intervals (see Table [1]). Interval 1 in black (asterisks), Interval 2 in green (diamonds), Interval 3 in red (triangles) and Interval 4 in blue (squares).
Fig. 5. The cyclotron line profiles of Fig. 3 for the four 35 d intervals (black crosses): Interval 1 (top left), Interval 2 (top right), Interval 3 (bottom left) and Interval 4 (bottom right), normalized to a common amplitude. The big red filled dots represent the weighted mean values of all four intervals. The solid blue line represents a best fit function (a combination of two cosine components). We conclude that, within the given statistics, there is no evidence for any change in shape of the profiles with 35 d phase.

background estimation files Sky\textsubscript{VLE} dated 2005 November 28 were used. We have investigated whether there are correlations between the fit parameters: no dependence of any single parameter on the others has been found. All spectral results stated below have been obtained through spectral fits with the assumption of negligible absorption due to neutral hydrogen, both intrinsic to the source and interstellar (which is at $5.1 \times 10^{19}$ cm\textsuperscript{-2}, Dal Fiume et al. 1998). The introduction of absorption as a free parameter does not lead to a reduction in $\chi^2$.

For the pulse phase resolved analysis, we have extracted the spectra in 14 pulse phase bins over the 1.24 s pulse, the widths of which were chosen to have roughly identical photon-statistics in each spectrum. The centers of the 14 phase bins are at pulse phases: 0.1, 0.3, 0.45, 0.55, 0.625, 0.675, 0.7125, 0.7375, 0.7625, 0.7875, 0.825, 0.875, 0.925, and 0.975. For the cyclotron line energy $E_{\text{cyc}}$ this is the finest resolution in pulse phase ever achieved (the four smallest bins around the peak of the pulse have a width of 1/80 of a phase).

With the spectral function described above, good fits with reduced $\chi^2$ between 0.9 and 1.2 are generally achieved. However, splitting up the data into small pulse phase intervals reveals that larger $\chi^2$ values appear: up to ~12 for the pulse phase range 0.775–0.85, and 1.4–2.8 for the pulse phase range 0.5–0.775. The main reason for this is the existence of the so called “10 keV feature” (see the discussion in 3.5). Modeling this feature by an extra Gaussian is generally successful and brings the $\chi^2$ down to acceptable values (in a few cases $\chi^2$ is still found around 1.3, which is most likely due to an imperfect modeling by a simple Gaussian).

3. Pulse phase variability

We concentrate on the detailed study of the variation of the following spectral parameters as function of pulse phase: the centroid energy $E_{\text{cyc}}$, the photon index $\Gamma$, and the intensity of the iron line at 6.4 keV. The phase profiles of these spectral parameters are compared with the shape of the flux profiles (the “pulse profiles”) for the four different 35 d phase intervals. In Figure 2 the 9–13 keV pulse profiles of the four intervals are shown. Note that the pulse profiles are of such high statistical quality that the uncertainties of the flux values (in each of the 128 pulse phase bins) are comparable to the width of the lines chosen to plot the profiles. The profiles are repeated in the following figures (as dashed lines) to allow a direct comparison with the profiles of the spectral parameters. For the remaining spectral parameters, such as the cutoff energy $E_{\text{cut}}$, the folding energy $E_{\text{fold}}$ (with reference to the used spectral model - highecut), and the width and depth of the cyclotron line ($\sigma_{\text{cyc}}$ and $\tau_{\text{cyc}}$), we will restrict ourselves to profiles of the combined data of cycle 323. For these parameters the statistics is such, that no statements can be made about a dependence on 35 d phase. In all Figures, the uncertainties given are one sigma (68%) values.

3.1. Cyclotron line energy

The profiles of the cyclotron line energy $E_{\text{cyc}}$ as a function of pulse phase are shown in Figures 3 and 4. Generally, $E_{\text{cyc}}$ roughly follows the shape of the pulse profile: the broad maximum is found to be close to the peak of the pulse profile (around pulse phase 0.7) for all 35 d phases. The formal uncertainties
Fig. 6. Photon-index profiles as function of pulse phase for Interval 1 (top left), Interval 2 (top right), Interval 3 (bottom left) and Interval 4 (bottom right).

of $E_{\text{cyc}}$ (as determined by XSPEC) are of the order of a few to several percent for pulse phases $0.0-0.6$ and $0.9-1.0$ and around one percent or less for pulse phases $0.6-0.9$ (close to the peak of the pulse). Adopting an additional systematic uncertainty of $0.15$ keV ($\sim 0.4\%$) for all data points, the profiles are well represented by cosine functions (see below). We find slightly different values for the mean $E_{\text{cyc}}$ and the peak-to-peak amplitude: a slight, but significant increase is found as function of pulse phase. The weighted mean values of the four highest values of $E_{\text{cyc}}$ of each group are $40.16 \pm 0.22$ keV, $40.89 \pm 0.20$ keV, $41.02 \pm 0.15$ keV and $41.52 \pm 0.20$ keV for mean $35$ d phases $0.03$, $0.10$, $0.15$ and $0.20$ (intervals 1 through 4), respectively. This corresponds to an increase by $\sim 0.7$ keV per $0.1$ units in $35$ d phase. The minimum $E_{\text{cyc}}$ values have larger uncertainties with a tendency to decrease with $35$ d phase, but, within uncertainties, they are consistent with a constant value (around $33.3$ keV). The peak-to-peak amplitude is then slightly increasing with $35$ d phase (with a mean of $\sim 7.6$ keV).

While the values for the maximum $E_{\text{cyc}}$ and the peak-to-peak amplitude are dependent on $35$ d phase, the sinusoidal shape of the modulation does not change. In order to verify this statement, the following analysis was performed: the original $E_{\text{cyc}}$ profiles (Figure 4) were scaled such that the separation of minimum to maximum $E_{\text{cyc}}$ had a constant value of $7.64$ keV, using the minimum values measured in phase bin $0.2-0.4$ and the weighted mean of the four highest values in each of the four profiles as reference. In this way, normalized $E_{\text{cyc}}$ profiles were generated. These normalized profiles are shown in Fig. 5 each in comparison to the weighted mean values and a smooth best fit function of the weighted mean: a combination of two cosine functions with common mean ($37.31$ keV), amplitude ($3.82$ keV), period and phase-zero values of $0.70$ and $0.772$ for pulse phases $0.3-0.875$ and $1.47$ and $0.3$ for the remaining phases, respectively. Visually, there is no obvious deviation of any of the four data sets from the weighted mean. As quantitative statistical test we performed the following: comparing the individual normalized profiles with that profile representing the weighted mean leads to $\chi^2$ values (normalized to the degrees of freedom) of $0.72$, $1.46$, $0.72$ and $1.32$ for the four intervals, respectively. As an
additional test, we have then performed a Monte Carlo simulation, in order to answer the following question. Assuming the profiles are indeed the same for all four intervals, that is equal to the best fit mean profile: what is the probability to find an individual profile which deviates from the mean profile in such a way that the $\chi^2$ reaches or exceeds the observed value? We have simulated 100 profiles for each of the four intervals, reproducing the statistical conditions of each interval (that is, the width of the Gaussian distribution from which the random numbers were drawn were chosen according to the uncertainty of the corresponding measured data point). Those probabilities are 68%, 29%, 73% and 4% for the four intervals, respectively. Even 4% for Interval 4 is a rather high and acceptable probability. We conclude, that, within our statistical accuracy, we find no evidence for any change in the shape of the $E_{\text{cyg}}$ profiles with 35 d phase.

### 3.2. Photon-index

Figure 6 shows the profiles of the photon index $\Gamma$ as function of the pulse phase. This spectral parameter is relatively constant for pulse phases up to 0.6 (showing a small positive slope), but has a distinct dip around the pulse peak and a smaller one around the left shoulder. The photon-indeces range between $0.42 \pm 0.01$ and $1.07 \pm 0.01$ for Interval 1, between $0.54 \pm 0.01$ and $1.08 \pm 0.02$ for Interval 2, between $0.36 \pm 0.01$ and $1.10 \pm 0.49$ for Interval 3 and between $0.33 \pm 0.01$ and $1.07 \pm 0.01$ for Interval 4. In general, the spectrum gets harder when the flux increases. While the main features of the pulse phase dependence of $\Gamma$ are very much the same in the four 35 d intervals, the depth of the new narrow dip around phase 0.62 shows a clear trend towards smaller values from Interval 1 to Interval 4 (see Fig. 7).

### 3.3. Iron line intensity

Another spectral parameter which shows a clear modulation as function of pulse phase is the iron line intensity (see Figure 8). It shows a moderate increase until the left shoulder of the pulse (phase $\sim 0.7$), after which it drops almost to zero around pulse phase 0.8, before reaching the start value again at pulse phase 0.9. The minimum is coincident with the peak of the X-ray pulse.

![Fig. 8](image-url). Iron line intensity profiles as function of pulse phase for Interval 1 (top left), Interval 2 (top right), Interval 3 (bottom left) and Interval 4 (bottom right). All intensities less than 1.0 are consistent with zero, therefore show two sigma upperlimits.

![Fig. 9](image-url). The iron line intensity profiles for the four different 35 d phase intervals (see Table 1): Interval 1 in black (asterisks), Interval 2 in green (diamonds), Interval 3 in red (triangles) and Interval 4 in blue (squares).
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3.4. Other spectral parameters

For the other spectral parameters $E_{\text{cut}}$, $E_{\text{fold}}$, the width $\sigma_{\text{cyc}}$, and the depth $\tau_{\text{cyc}}$, we show their pulse phase dependence only for the total data sets (not resolved in 35 d phase), because in these parameters we find a considerably larger scatter, such that no conclusions can be drawn about a possible variation (or not) with 35 d phase (the 35 d phase resolved profiles can be found in Vasco 2012). The profiles are shown in Figures 10. Also for these parameters there is significant modulation with pulse phase, and they generally show a broad maximum around the phase of the pulse peak, with the interesting feature that for $E_{\text{fold}}$ there is a dip in this maximum, right at the pulse peak. It is also worth noting, that, except for $E_{\text{fold}}$, the measured values for pulse phases $<0.4$ are rather high, not following the shape of the pulse profiles.

3.5. The “10 keV feature”

It has been noticed (e.g. Coburn et al. 2002), that in fitting observed spectra of accreting X-ray binaries with a cut-off power law a “wiggle” or “bump” around $\sim 10$ keV to $\sim 15$ keV is found in the residuals. This applies to many different sources including Her X-1 (also to non-cyclotron line sources), to observations with different satellites (e.g. Ginga, Beppo/SAX and RXTE), and to analysis with different spectral functions involving a power law (e.g., highcut, plcut, Fermi-Dirac or NPEX). The feature is therefore considered to be inherent to accreting X-ray binaries. So far, no physical explanation has been suggested.

In our current analysis, using data from observations of Her X-1 by RXTE, this feature does not play a role in the phase averaged analysis. However, in the pulse phase resolved analysis, larger $\chi^2$ values appear due to the existence of this ”10 keV feature”: up to $\sim$12 for the pulse phase range 0.775–0.85, and 1.4–2.8 for the pulse phase range 0.5–0.775. Modeling this feature by an extra Gaussian is generally successful and brings the $\chi^2$ down to acceptable values (in a few cases $\chi^2$ is still found around 1.3, which is most likely due to an imperfect modeling by a simple Gaussian). The centroid energy is generally found at $\sim 16$ keV with a $\sigma$ of $\sim 5$ keV. Fig. 11 shows an example of the spectrum of Interval 3, pulse phase bin 0.8–0.85, where the $\chi^2$ is reduced from 5.55 to 1.38 for 224 dof. We have verified, that all the other spectral parameters are not altered when introducing this extra Gaussian. We note that this feature occurs most strongly at a pulse phase close to the peak of the pulse, where all other spectral parameters also tend to show their strongest variability.

4. Summary and discussion

We have selected the best data available in the RXTE archive of an observation of a Main-On of Her X-1 to perform detailed pulse phase resolved spectroscopy in the energy range 3.5–75 keV. The observations of November 2002 (cycle no. 323) provide the best RXTE coverage of a Main-On (Vasco 2012), such that an attempt could be made to study pulse phase dependent spectral variations also as a function of the phase of the 35 d modulations - flux and shape of the pulse profiles -, which are believed to be due to precessional motion of the accretion disk.

(see Figure 8 and 9) in a similar way as the photon index. The iron line intensity ranges between 0 (for all intervals) and 12.40 for Interval 1, 14.62 for Interval 2, 13.19 for Interval 3 and 10.75 for Interval 4. The profiles of 35 d intervals 2–4 are quite similar to one another, while for Interval 1, the iron line intensity is significantly lower for pulse phases $<0.6$.

Fig. 10. Pulse phase profiles for the cutoff energy $E_{\text{cut}}$ (top left), the folding energy $E_{\text{fold}}$ (top right), the width of the cyclotron line $\sigma_{\text{cyc}}$ (bottom left), and the depth $\tau_{\text{cyc}}$ of the cyclotron line (bottom right) for the total data sets.

6 http://tobias-lib.uni-tuebingen.de/frontdoor.php?source=opus=6346
and, possibly, the neutron star itself. As spectral model we have used a power law with an exponential cut-off (the hihecut model in XSPEC) modified by two line features, a multiplicative cyclotron line (in absorption) and an additive iron fluorescence line at 6.4 keV. For these three parameters, the analysis was performed for four one day long intervals at different mean 35 d phases (0.03, 0.10, 0.15 and 0.20, respectively). For the other spectral parameters of the applied spectral model, no 35 day phase dependence was studied.

Fig. 11. Residuals of the X-ray spectrum of pulse phase bin 0.80 - 0.85 for Interval 3 without (top) and with (bottom) an extra Gaussian component in the fit.

4.1. Pulse phase dependence

The results of the pulse phase resolved analysis can be summarized as follows. We distinguish between results which confirm and improve in statistical significance which was known before and new results.

4.1.1. Confirmed and improved results

- All spectral parameters show a strong modulation with pulse phase.
- The strongest variations always appear around the peak of the pulse (at maximum flux).
- The profile of the centroid energy $E_{\text{cy}}$ of the cyclotron line follows roughly the modulation of the pulse profile (maximum energy around maximum flux). The shape is close to "sinusoidal".
- The photon index profile (Fig. 6 and 7) shows a strong dip close to the main peak. The dip around the main peak is a known feature, nicely confirmed here and very well resolved in this analysis.

4.1.2. New results

- It is the first time that for Her X-1 a phase profile of the centroid cyclotron line energy $E_{\text{cy}}$ is produced with such a high resolution and statistical accuracy. The four smallest bins around the peak of the pulse have a width of 1/80 of a phase. The mean variation (minimum to maximum) is $\sim 23\%$ (Figs. 3 and 4). The modulation is close to "sinusoidal"; however, at least two sinusoidal components (with greatly different periods) are needed to model the modulation.

- For 35 d intervals 1–3, the peak of the $E_{\text{cy}}$ profile is structured (see Fig. 3), with local maxima near the maximum of the main peak of the pulse profile and around the position of the left shoulder of the main peak. A similar structure is seen in $\tau_{\text{cy}}$.
- The folding energy $E_{\text{fold}}$ follows the pulse profile, except for a narrow dip in the peak of the pulse (Fig. 10 upper-right).
- Similarly, the cut-off energy $E_{\text{cut}}$, the width and depth of the cyclotron line ($\delta_{\text{cy}}$ and $\tau_{\text{cy}}$), roughly follow the pulse profile, except for higher values at phases 0.0–0.5 (Fig. 10). While the strong dip in the photon index profile (Fig. 6 and 7) close to the main pulse peak was known (even though less well measured), the second, very narrow dip (around phase 0.61) is a newly detected feature. We note that this is close to the position of the left shoulder component of the main peak.
- The profile of the intensity of the 6.4 keV iron line shows a distinct minimum at maximum flux (Figs. 8 and 9).
- We find that the so called "10 keV feature" seen in the spectra of several accreting binary pulsars, is around 16 keV in Her X-1, with a strong pulse phase dependence: essentially concentrated in the narrow pulse phase interval 0.8–0.9. Spectral fits for all the other phase intervals do not show significant residuals.

4.2. Precession phase dependence

With respect to any variation of the spectral parameters (or their pulse phase variations) with 35 d phase, we consider only the cyclotron line energy $E_{\text{cy}}$, the photon index $\Gamma$ and the intensity of the iron fluorescence line at 6.4 keV. For the other parameters the photon statistics is not sufficient to distinguish between different 35 d phases. It is the first time that a study of the variation of spectral parameters with 35 d phase was done. So, the following results are all new.

- The maximum value for $E_{\text{cy}}$ (reached near the pulse peak) shows a slight, but significant increase with 35 d phase. The mean value is $41.0 \pm 0.1$ keV, and the slope is $\sim 0.7$ keV per 0.1 units of 35 d phase.
- The minimum value for $E_{\text{cy}}$ (around pulse phase 0.3) is formally consistent with a constant value of $33.5 \pm 0.5$ keV (we would not rule out a small decrease). Together with the maximum increasing, this means that also the peak-to-peak amplitude (mean value $7.6 \pm 0.5$ keV) is increasing with 35 d phase.
- The shape of the $E_{\text{cy}}$ profile (see Figs. 4 and 5) is independent of 35 d phase.
- Similarly, the shape of the $\Gamma$ profiles is independent of 35 d phase, except for the new narrow dip around pulse phase 0.62 (Figs. 6 and 7), where a trend to a smaller depth from Interval 1 to Interval 4 is seen.
- The shape of the iron line intensity profiles are similar to one another, except for Interval 1 which shows a distinctly lower intensity for pulse phases <0.6 (Figs. 8 and 9).

4.3. Comparison with previous results and physical conclusions

For Her X-1, the dependence of the spectral parameters on pulse phase have been studied by several authors on the basis of observations with different X-ray instruments (see Introduction). In the following we highlight our new findings and significant improvements and discuss possible physical implications.
Pulse phase dependence of the centroid cyclotron line energy: This modulation is quite common among accreting X-ray pulsars and is generally believed to be due to the changing viewing angle under which the X-ray emitting regions are seen (e.g., Kreykenbohm et al. 2004 and references therein). If we adopt the idea, that peaks (high X-ray flux) in the pulse profile can be associated with beamed radiation emitted from the accreting regions around the magnetic poles, then the observed pulse profiles are a representation of the emission characteristics. During a full neutron star rotation we may be seeing different emitting spots (most likely from a complex magnetic field configuration which deviates from a simple dipole (see e.g. Suchy et al. 2008), possibly having a multi-pole configuration) and the emerging beamed radiation under changing angles. The situation is further complicated by gravitational bending. It is therefore not straightforward to interpret observed pulse profiles in terms of accretion geometry and beaming characteristics (see e.g. Caballero et al. 2011).

If we, however, assume that the main peak in the pulse profile of Her X-1 is due to a reasonably narrow pencil beam (following e.g. Pravdo et al. 1978; Klochkov et al. 2008), emitted perpendicular to the neutron star surface at the hot spot of the main accreting material, then we may associate this pulse phase with a situation where the observer is looking down into the accretion column, parallel to the magnetic field lines and the velocity of the in-falling material. At this phase we see the maximum cyclotron line energy and the hardest spectrum (a sharp minimum in photon index $\Gamma$) - that is, we are looking down to small distances from the neutron star surface where the magnetic field is strongest and the temperature is highest. Here, we also find the largest width and depth of the cyclotron line, as well as the highest $E_{\text{cut}}$ and a local dip within a broad peak of $E_{\text{cyc}}$.

Along similar lines, the newly discovered structure of a second (very sharp) minimum in photon index $\Gamma$ around pulse phase 0.61 (Fig. 2), corresponding to the left shoulder of the main peak, may be the signature of a second pole. At this pulse phase we also find indications for another relative maximum in $E_{\text{cyc}}$ (Fig. 3).

Pulse phase dependence of the 6.4 keV iron line flux: Figs. 8 and 9 show a sharp and deep minimum of the intensity of the 6.4 keV iron line, again coincident with the main pulse peak. The profile is very similar to that of $\Gamma$ (Fig 2) and the Fe line intensity drops to near zero at pulse phase $\sim 0.79$ (where $\Gamma$ is minimal). It is the first time, that such a sharp minimum in the 6.4 keV flux is observed. While Choi et al. (1994), in observations by Ginga, and Zane et al. (2004), in observations by XMM-Newton, had found a sinusoidal-like modulation with a broad minimum around the pulse peak, Oosterbroek et al. (2000), with data from BeppoSAX, were only able to measure the Fe-I line around 1 keV, for which they had found a minimum around the pulse peak which was similarly sharp and deep as we now find for the 6.4 keV line. It is generally believed, that the origin for the Fe lines is fluorescence in relatively cool material which is being illuminated by the X-ray beam. For the geometrical site, however, where this fluorescence is taking place, or whether it is predominately in reflection or in transmission, no general consensus has been reached so far. Several different regions have been proposed (see e.g. Choi et al. 1994 and references therein): 1) the outer edge of the accretion disk extended in the vertical direction; 2) a corona of the accretion disk; 3) the atmosphere of the companion star; 4) a surface at the Alfven radius. The Alfven surface is currently the most popular one. However, also the accretion column was proposed (e.g. Leahy 2001 and Inam & Baykal 2005), despite the argument by White et al. (1983) that the material in the column may be too hot.

We like to stress the observed fact, that at the pulse phase where we see the highest continuum flux there is essentially no Fe line flux. This should provide some constraints in finding a feasible geometry. If the above adopted scenario is correct, by which we look down into the pencil beam at the pulse peak, the Fe line profile may be consistent with a hollow cone geometry, where the majority of the radiation escapes through the empty center of the cone, where there is no material which can fluorescence, while the walls are seen at off-center angles. It is, of course, assumed that the bottom of the cone close to the neutron star surface is filled with X-ray emitting plasma.

The non-dependence of the cyclotron line energy profile on 35 d phase and the model of neutron star free precession: One of the results of this study is the finding that the shape of the cyclotron line energy profile ($E_{\text{cyc}}$ vs pulse phase) is independent of 35 d phase (Figs. 4 and 5), while the maximum cyclotron line energy increases slightly by $\sim 0.7$ keV per 0.1 units of 35 d phase. If we attribute the strong variations in pulse profile with 35 d phase to precession of the neutron star, that is to the change of the angle under which we see the X-ray emitting regions, we would expect also the shape of the cyclotron line energy profile to change with 35 d phase. However, this expectation needs a quantitative analysis. Recently concluded detailed model calculations (Postnov et al. 2012), assuming neutron star free precession, have shown, that under a specific set of assumptions (including a certain multi-pole geometry and emission into narrow pencil beams), the observed pulse profiles in the energy range 9-13 keV (the same as used in Staubert et al. 2010, 2013) can be reproduced with high accuracy. We point out here that the model of Postnov et al. (2012) is an attempt to give a physical explanation of the variation in pulse profiles by neutron star free precession, requiring several specific assumptions, while the template approach of Staubert et al. (2013), mentioned in the Introduction, does not need any physical assumptions, it just makes use of the observed pulse profiles. We plan to investigate whether the here described constancy of the shape of the cyclotron energy profile is consistent with the physical model. In this sense our result opens a new channel to test the model of free precession of the neutron star.

References

Boynton, P. E., Crosa, L. M., & Deeter, J. E. 1980, ApJ, 237, 169
Caballero, I., Kraus, U., Santangelo, A.; Sasaki, M., & Kretschmar, P. 2011, A&A, 526, A131
Choi, C. S., Nagase, F., Makino, F., et al. 1994, ApJ, 437, 449
Coburn, W., Heindl, W. A., Rothschild, R. E., et al. 2002, ApJ, 580, 394
Dal Fiume, D., Orlandini, M., Cusumano, G., et al. 1998, A&A, 329, L41
Deeter, J. E., Scott, D. M., Boynton, P. E., et al. 1998, ApJ, 502, 802
Endo, T., Nagase, F., & Mihara, T. 2000, PASJ, 52, 223
Ferrigno, C., Falanga, M., Bozzo, E., et al., 2011, in The X-ray Universe 2011, ed. J-U. Ness & M. Eilek, 69–
Gerend, D. & Boynton, P. E. 1976, ApJ, 209, 562
Giacconi, R., Gursky, H., Kellogg, E., et al. 1973, ApJ, 184, 227
Inam, S. C. & Baykal, A. 2005, MNras, 361, 139
Kahabka, P. 1987, MPE Rep., No. 204, 139 pp., 204
Klochkov, D., Shakura, N., Postnov, K., Staubert, R., & Wilms, J. 2006, ArXiv Astrophysics e-prints
Klochkov, D., Staubert, R., Postnov, K., et al. 2008, A&A, 482, 907
Kreykenbohm, I., Wilms, J., Coburn, W., et al. 2004, A&A, 427, 975
Kunz, M., Kendziorra, E., Kretschmar, P., et al. 1996, A&AS, 120, C233
Leahy, D. A. 2001, in International Cosmic Ray Conference. Vol. 6, International Cosmic Ray Conference, 2536
Lutovinov, A. A., Grebenev, S. A., Pavlinsky, M. N., & Sunyaev, R. A. 2000, Astronomy Letters, 26, 691
McCray, R. A., Shull, J. M., Boynton, P. E., et al. 1982, ApJ, 262, 301
Oosterbroek, T., Parmar, A. N., Dal Fiume, D., et al. 2000, A&A, 353, 575
Postnov, K., Shakura, N., & Staubert, R. 2012, MNRAS, submitted
Pravdo, S. H., Bussard, R. W., Becker, R. H., et al. 1978, ApJ, 225, 988
Pravdo, S. H., Bussard, R. W., & White, N. E. 1979, MNRAS, 188, 5P
Reynolds, A. P., Quaintrell, H., Still, M. D., et al. 1997, MNRAS, 288, 43
Rothschild, R. E., Markowitz, A., Rivers, E., et al. 2011, ApJ, 733, 23
Scott, D. M. & Leahy, D. A. 1999, ApJ, 510, 974
Scott, D. M., Leahy, D. A., & Wilson, R. B. 2000, ApJ, 539, 392
Soong, Y., Gruber, D. E., Peterson, L. E., & Rothschild, R. E. 1990, ApJ, 348, 641
Staubert, R., Klochkov, D., Vasco, D., et al. 2012, A&A, 00, 00, accepted 05 Dec 12, astro-ph.HE/1212.5472
Staubert, R., Klochkov, D., Vasco, D., & Wilms, J. 2010a, in Proceedings of the 8th INTEGRAL Workshop "The Restless Gamma-ray Universe" (INTEGRAL 2010). September 27-30 2010. Dublin Castle, Dublin, Ireland. Published online at http://pos.sissa.it/cgi-bin/reader/conf.cgi?confid=115, arXiv: astro-ph.HE 1110.6717
Staubert, R., Klochkov, D., Vasco, D., & Wilms, J. 2010b, in Proceedings of the 8th INTEGRAL Workshop "The Restless Gamma-ray Universe" (INTEGRAL 2010). September 27-30 2010. Dublin Castle, Dublin, Ireland. Published online at http://pos.sissa.it/cgi-bin/reader/conf.cgi?confid=115
Staubert, R., Klochkov, D., & Wilms, J. 2009, A&A, 500, 883
Staubert, R., Shakura, N. I., Postnov, K., et al. 2007, A&A, 465, L25
Suchy, S., Pottschmidt, K., Wilms, J., et al. 2008, ApJ, 675, 1487
Tananbaum, H., Gursky, H., Kellogg, E. M., et al. 1972, ApJ, 174, L143+
Trümper, J., Kahabka, P., Oegelman, H., Pietsch, W., & Voges, W. 1986, ApJ, 300, L63
Trümper, J., Pietsch, W., Reppin, C., et al. 1978, ApJ, 219, L105
Vasco, D. 2012, PhD thesis, Univ. of Tübingen, Germany
Vasco, D., Klochkov, D., & Staubert, R. 2011, A&A, 532, A99+
Voges, W., Pietsch, W., Reppin, C., et al. 1982, ApJ, 263, 803
White, N. E., Swank, J. H., & Holt, S. S. 1983, ApJ, 270, 711
Zane, S. & Ramsay, G. 2001, in Two Years of Science with Chandra
Zane, S., Ramsay, G., Jimenez-Gurate, M. A., Willem den Herder, J., & Hailey, C. J. 2004, MNRAS, 350, 506