SIR THOMAS WALTER BANNERMANN KIBBLE
23 December 1932 — 2 June 2016
Professor Tom Kibble was an internationally-renowned theoretical physicist whose contributions to theoretical physics range from the theory of elementary particles to modern early-Universe cosmology. The unifying theme behind all his work is the theory of non-abelian gauge theories, the Yang–Mills extension of electromagnetism. One of Kibble’s most important pieces of work in this area was his study of the symmetry-breaking mechanism whereby the force-carrying vector particles in the theory can acquire a mass accompanied by the appearance of a massive scalar boson. This idea, put forward independently by Brout and Englert, by Higgs, and by Guralnik, Hagen and Kibble in 1964, and generalized by Kibble in 1967, lies at the heart of the Standard Model and all modern unified theories of fundamental particles. It was vindicated in 2012 by the discovery of the Higgs boson at CERN. According to Nobel Laureate Steven Weinberg: ‘Tom Kibble showed us why light is massless’; this is the fundamental basis of electromagnetism.

1. Early years

1.1. India

Thomas Walter Bannerman Kibble was born on 23 December 1932 in Madras (now Chennai), Madras Presidency, India. Both his parents were from missionary backgrounds. His father, Fred Kibble, was a professor of mathematics at Madras Christian College. His mother, Janet Bannerman, descended from a long line of ministers in the Church of Scotland. Her father,
William Burney Bannerman, was a doctor in the Indian Medical Service. Tom spent his first 10 years in India, and was educated in the Doveton Corrie School, rated as one of the best schools in Madras. In the summers, his family generally went to Kodaikanal in the Palni Hills to escape the heat. He particularly enjoyed rowing on the lake there (figure 1).

The principal mathematical influence on young Tom was his father, who was particularly interested in the symmetry patterns of Moghul architecture in Agra and elsewhere in India. From this exposure, he gained a love of geometry. His father taught him about the regular and semi-regular solids, and he enjoyed making paper models of them.

It had always been intended that at about the age of 10, Tom should return to Britain for later schooling. However, the timing fell during the Second World War, and travel was difficult. In early 1943, passage became available at just two days’ notice—he was not even able to say
goodbye to his friends. Tom went on this voyage alone. Since the Suez Canal was closed, the voyage was around Africa, then across the Atlantic and back in convoy, taking 11 weeks. In Capetown, Tom went ashore alone at one point, but upon returning to the ship found that it was gone; happily, it had simply moved to another dock. Such experiences gave young Tom a self-reliance which remained with him for life.

### 1.2. Edinburgh

From 1944 to 1951, Tom attended Melville College, Edinburgh, becoming head boy in 1950–1951 (figure 2). In his final years at school, he became very interested in cars and, together with friends, owned a series of old wrecks. His father was keen for him to go to Cambridge. However, in his recollections Tom said that he found the problems for the scholarship exams extremely tedious (and he added that by then he was enjoying life too much to work on them), so he was not successful. Consequently, he went to the University of Edinburgh, a decision that he never regretted.

Two important influences on Tom’s future career while an undergraduate in Edinburgh were Robin Schlapp and Andrew Nisbet, who taught classical mathematical physics. Also important was the teaching in pure mathematics in the mathematics department under Alexander Aitken FRS—who was known for his great facility for mental arithmetic, being able to extract cube roots of twelve-digit numbers in seconds. Tom remarked that it was no accident that a considerable number of the leading mathematical physicists in the UK were educated in Edinburgh during this period. Two other near contemporaries in Edinburgh were David Olive (FRS 1987) and Keith Moffat (FRS 1986).

Halfway through Tom’s undergraduate course, Max Born FRS retired and was replaced by Nicholas Kemmer (FRS 1956), who gave inspiring lectures completely without notes (but with the distracting habit of balancing a row of pieces of chalk end on end). Tom continued on to a PhD in Edinburgh—his supervisor was John Polkinghorne (FRS 1974)—and his PhD thesis consisted of two parts: Schwinger’s action principle; and dispersion relations for inelastic scattering processes.

During Tom’s education in Edinburgh, he was for some time a member of the Student’s Representative Council, and it was through this that he met Anne Allan, who was also a member (see figure 3). They became engaged in June 1954 and married when he was in the middle of his PhD, in July 1957. They had three children: Helen, Alison and Robert.

### 1.3. Caltech

Following the award of his PhD, Tom received a Harkness Fellowship to spend a year in the United States. Tom and Anne went to Caltech, where Tom got to know Dick Feynman (ForMemRS 1965) and Murray Gell-Mann (ForMemRS 1978). Upon returning to the UK in 1959 as a NATO Fellow, Tom joined the Theoretical Physics Group at Imperial College.

### 2. Imperial College London

#### 2.1. Quantum field theory and the influence of Abdus Salam

In 1957 the head of the physics department at Imperial, P. M. S. Blackett FRS, later Lord Blackett (PRS 1965–1970), recruited the Pakistani physicist Abdus Salam (FRS 1959) from Cambridge to be founding member of the Theoretical High Energy Physics Group. Salam
remained as professor there until his death in 1996, sharing the 1979 Nobel Prize with Sheldon Glashow and Steven Weinberg (ForMemRS 1981) for unifying the weak and electromagnetic forces (figure 4). As Tom himself recalls (19)*: ‘It was a very exciting place to be. We had lots of visitors: Steven Weinberg, Murray Gell-Mann, Ken Johnson, Lowell Brown (3), Stanley Mandelstam, John Ward (FRS 1965), to name but a few. The year I arrived, 1959, was also the year Salam became the youngest Fellow of the Royal Society at the age of 33.’

* Numbers in this form refer to the bibliography at the end of the text.
In the late 1950s and early 1960s, however, the prevailing mood in fundamental physics was not favourable to quantum field theories other than quantum electrodynamics (QED). The success of QED, pioneered by Richard Feynman, Julian Schwinger, Shinichiro Tomanaga and Freeman Dyson (FRS 1952), relied on the smallness of the coupling constant and hence a sensible perturbative expansion. This would not be the case for the strong interactions in the Yukawa theory, however, as the squared coupling constant would be about 15. In Tom’s own words (19): ‘Salam was convinced from an early date—certainly well before
he came to Imperial College—that the fundamental theory would be a gauge theory. This was not a popular view. Many people thought that field theory had had its day, particularly for the strong interaction, where it was to be superseded by S-matrix theory. This was an appealingly economical idea: everything was supposed to follow from fundamental principles of covariance, causality and unitarity. But Salam was undeterred and under his influence, and that of P. T. Matthews (FRS 1963), the construction of unified gauge theories was a major theme of the work at Imperial. Tom’s interest in gauge field theories extended throughout his career. In 1961, he wrote an influential paper on the relation between gravity and a gauge theory of the Poincaré group (1, 14). Another significant contribution was a series of four papers (5–8) devoted to resolving the notoriously tricky problem of defining an S-matrix in theories with massless particles such as QED. There has been a recent revival of interest in this problem, and Tom’s resolution, involving coherent states, continues to be frequently cited (see, for example, Hannesdottir & Schwartz 2020).

Tom concludes (19): ‘I shall always look back on these early years as the most exciting of my life. There was a wonderful atmosphere at the College; I cannot think of anywhere I would rather have been. I shall always be immensely grateful to Abdus Salam for making it possible for me to participate in these memorable developments.’ The inspiration was reciprocated: Salam’s 1968 electroweak unification relied crucially on the spontaneous-symmetry-breaking mechanism that Tom pioneered (2, 4).
2.2. Spontaneous symmetry breaking

The first example of a gauge theory beyond QED was the $SU(2)$ non-abelian gauge theory proposed by C.-N. Yang (ForMemRS 1992) and R. L. Mills (Yang & Mills 1954) and by Salam’s student Shaw (1955). Although originally intended for the strong interactions, this formed the basis of Salam’s unification of the weak and electromagnetic forces. One of the key problems for the weak interactions was their short range, which implies that the force carrier particles must be massive, as opposed to the naturally massless excitations of gauge theories.

In 1964–1965, Gerry Guralnik and Dick Hagen were visitors. This led to a key collaboration with Tom that borrowed ideas from P. Anderson (ForMemRS 1980) in condensed matter physics in the context of the Ginzburg–Landau model of superconductivity. This early model (now derivable from the full Bardeen–Cooper–Schrieffer theory as an effective theory) involves a scalar order parameter field $\phi(t, x)$ representing the wave function of a condensate of Cooper pairs, i.e. bound state pairs of electrons. The Hamiltonian for the $\phi$ field is

$$ H = \int d^3x \left[ \frac{1}{2m} \mathbf{D}\phi^* \cdot \mathbf{D}\phi + V(\phi) \right], $$

where $\mathbf{D}\phi = \nabla\phi - 2ie\mathbf{A}\phi$ is the electromagnetic covariant derivative. Expanding in a power series, $V(\phi) = \alpha\phi^*\phi + \frac{1}{2}\beta(\phi^*\phi)^2$ when $\phi$ is small. Such a potential has a ‘sombrero’ shape, shown in figure 5, if $\alpha < 0$, $\beta > 0$.

The minima of $V(\phi)$ do not occur at the symmetrical point $\phi = 0$ but rather at points in the ‘rim’ where $\phi^*\phi = -\frac{\alpha}{\beta}$. This shift to a minimum in the ‘rim’ also produces a term $-\frac{2ae^2}{\beta M} \mathbf{A} \cdot \mathbf{A}$ giving a mass to the electromagnetic vector field $\mathbf{A}$. Consequently, the photon becomes massive in the broken-symmetry phase of a superconductor (Anderson 1963).

The idea that a similar mechanism could also give masses to elementary fermions in a relativistic field theory had been suggested by Nambu (1960). However, the Ginzburg–Landau model involves an effective scalar-field order parameter that is clearly not a fundamental field of the superconducting theory. The model was also non-relativistic, and the complications
imposed by relativity seemed at the time to be crucial. Moreover, giving a mass to the photon was something definitely not wanted in a fundamental relativistic theory.

### 2.3. Goldstone modes

Nambu and Jona-Lasinio proposed in 1961 a four-fermion interaction model based on a spinor field $\psi$ with both an ordinary $U(1)$ rigid phase symmetry $\psi \to e^{i\alpha} \psi$ and also a rigid chiral $U(1)$ phase symmetry $\psi \to e^{i\tilde{\alpha}\gamma_5} \psi$ with a Lagrangian interaction term $L_{\text{int}} = g\left[ (\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\psi)^2 \right]$ (Nambu & Jona-Lasinio 1961a, b). In this model, a condensate of the $\psi$ field occurs: $\langle \bar{\psi}\psi \rangle \neq 0$. The model successfully gives a mass to the fermion, but it also gives rise to a massless pseudoscalar mode. Nambu and Jona-Lasinio proposed that, given additional soft explicit chiral symmetry breaking, the pseudoscalar mode could in the end have a small mass and could be identified with the pion. (Indeed, such explicit chiral symmetry breaking is understood now to arise from quark masses in the underlying QCD theory.)

The fact that spontaneous symmetry breaking of rigid symmetries inevitably leads to the existence of massless modes is the content of Goldstone’s theorem, named after Jeffrey Goldstone (FRS 1974), of which a proof was given in 1962 by Goldstone, Salam and Weinberg (Goldstone et al. 1962). This is illustrated by the two modes of a complex scalar sketched in figure 5. The radial mode clearly experiences a restoring force, but the tangential ‘rolling mode’ does not. In a relativistic theory, the radial mode experiencing a restoring force corresponds to a massive particle, while the rolling mode corresponds to a massless particle.

The seeming inevitability of Goldstone modes arising in consequence of symmetry breaking posed a real obstacle to the implementation of the symmetry-breaking idea in relativistic theories, since no such massless bosons were known in physics. Attempts to circumvent the occurrence of such massless Goldstone modes by the addition of explicit symmetry-breaking mass terms were not acceptable when considered at the level of fundamental theory because such terms give rise to non-renormalizability, i.e. uncontrollable behaviour at high energies. Such behaviour could be contemplated in an approximate effective theory (such as in pion physics) but not at the level of a fundamental theory of microphysics.

### 2.4. The Brout–Englert–Higgs–Guralnik–Hagen–Kibble mechanism

The resolution of the Goldstone-mode problem came within a short period in 1964 from the work of three independent groups: Brout and Englert; Higgs; Guralnik, Hagen and Kibble (BEHGHK).

The key point was the recognition that Goldstone’s theorem could be evaded in a gauge theory. Peter Higgs (FRS 1980) had previously shown that, for a gauge theory formulated in radiation gauge, the lack of explicit relativistic Lorentz invariance arising from the gauge choice renders Goldstone’s theorem inapplicable.

The three groups approached the symmetry-breaking problem from different perspectives:

- Brout and Englert, in the first paper to be published (Brout & Engelert 1964), based their argument on a calculation of vacuum polarization in lowest-order quantum perturbation theory about the symmetry-breaking vacuum state.
- Higgs’s treatment was purely classical and was similar to the Ginzburg–Landau model of superconductivity (which he did not know at the time), but in a relativistic model (Higgs 1964).
Guralnik, Hagen and Kibble used a quantum operator formalism, concentrating on the role of the current conservation law and studying the precise way in which the Goldstone theorem could be evaded (2).

The key conclusion of this fundamental development was that, in a gauge theory, the Goldstone mode can be absorbed into the longitudinal mode of the vector gauge field, rendering it massive. The existence of a residual scalar mode, which was discussed only in Higgs’s paper, was not considered important at the time.

2.5. Tom Kibble’s 1967 paper, the massless photon and the Standard Model

The three BEHGHK 1964 papers addressed the key problem of how to avoid Goldstone’s theorem and generate masses for vector fields by spontaneous symmetry breaking. It was hoped that this could lead to a renormalizable theory with controlled ultraviolet behaviour. (This was later shown by Gerard ’t Hooft and Martinus Veltman (’t Hooft & Veltman 1972).) However, another key issue remained for a successful application to particle physics: how can one ensure that some gauge fields remain massless, and in particular the photon field?

The three 1964 symmetry-breaking papers had attracted little attention at the time. Meanwhile, Glashow (1961) and Salam and Ward (unaware of Glashow’s work) (Salam & Ward 1964) proposed a unified model of the weak and electromagnetic interactions based on the gauge group $SU(2) \times U(1)$, but without any explanation for the generation of mass needed for the vector particles responsible for the weak interactions. What was needed to pull these diverse elements together to form what we now know as the Standard Model was a deeper understanding of the general structure of spontaneous symmetry breaking in gauge theories. In 1967, Tom wrote a magisterial paper (4) laying out the complete mathematical structure of spontaneous symmetry breaking and the Higgs effect. It provided the complete explanation of which vector gauge fields become massive and which remain massless after spontaneous symmetry breaking and the Higgs effect.

This paper was acknowledged by Steven Weinberg in his 1967 paper as settling the issue of which gauge fields remain massless (Weinberg 1967). This was crucial to the creation of a unified model of the weak and electromagnetic interactions, as the quantum of the electromagnetic field, the photon, had to remain massless. This model, stemming from the work by Sheldon Glashow (Glashow 1961) and culminating in the 1967 and 1968 papers of Weinberg (Weinberg 1967) and Salam (Salam 1968), led to the award of the Nobel Prize for Physics in 1979 (see figure 4).

A major part of what is now known as the Standard Model, this unification of weak and electromagnetic interactions is based upon the electroweak gauge group $SU(2) \times U(1)$, which has four Lie algebra generators. The corresponding four vector gauge fields couple to a Higgs complex scalar-field gauge-symmetry doublet with a Lagrangian potential term similar to the ‘sombrero’ potential sketched in figure 5. The Higgs complex doublet contains four real fields. Three of these belong to the broken-generator coset $(SU(2) \times U(1))/U(1)_{\text{em}}$, where the unbroken electromagnetic subgroup $U(1)_{\text{em}}$ mixes one combination of $SU(2)$ generators with the $U(1)$ ‘hypercharge’ factor of the original gauge group prior to symmetry breaking. These erstwhile Goldstone fields are absorbed by the Higgs effect into masses for three of the vector gauge fields. The gauge field for the fourth gauge group generator, corresponding
to the unbroken $U(1)_{\text{em}}$ subgroup, remains massless. This is the electromagnetic photon field $A_\mu$.

To date, the electroweak Standard Model is one of the most precisely tested and verified achievements of elementary particle physics.

### 2.6. Cosmic strings

In 1976, Tom once more wrote an extremely influential paper on the consequences of spontaneous symmetry breaking, entitled ‘The topology of cosmic domains and strings’ (9). The focus now was on the topological structure of solutions and domain formation during phase transitions.

When applied to proposed extensions of the unification idea to include also the strong interactions, spontaneous symmetry breaking leads to non-singular magnetic ‘monopole’ solutions, as found by Gerard ‘t Hooft (‘t Hooft 1974) and Alexander Polyakov (Polyakov 1974) in an $SU(2)$ model with a triplet Higgs field. These involve non-trivial maps from the 2-sphere at spatial infinity to the vacuum manifold $SU(2)/U(1)$. Tom’s 1976 paper (9) considered the cosmological implications of more general ‘topological defects’: topologically stabilized strings and domain walls as well as monopoles. These possibilities are classified by the homotopy groups $\pi_k(M)$ of the Higgs vacuum manifold, $M$, i.e. by maps from an $S^k$ sphere at spatial infinity to $M$. If $\pi_0(M) \neq 0$, one can have domain walls; if $\pi_1(M) \neq 0$, one can have string solutions; and if $\pi_2(M) \neq 0$, one can have monopoles. Tom realized that these structures could condense as the Universe cooled from the hot conditions prevailing in the Big Bang, and might therefore have striking effects on the development of large-scale structure in the Universe.

Finite-temperature corrections induce changes to the effective potential for the Higgs scalar fields, with an unbroken gauge symmetry $G$ being restored above some critical temperature $T_c$. Below this temperature, phase transitions will take place, producing a phase domain structure separated by such topological defects. The values of the symmetry-breaking order parameter can be roughly constant in widely separated domains, but can be in conflict where domains meet. This causes formation of a topological defect, in the centre of which the order parameter rises to its unbroken symmetry value, with a consequent concentration of energy. In the Big Bang origin of the universe, gauge symmetries would be restored by the initial high temperatures, with a rapid thermal freeze-out quench as the temperature decreased below $T_c$.

The implications of $\pi_0(M) \neq 0$ domain walls are very bad for cosmology, since they would be very massive and would cause an impossibly large anisotropy in the 3 K cosmic microwave background (9). Consequently, the Higgs vacuum manifold must be connected. This constraint on models of particle physics heralded a very fertile interaction between cosmology and fundamental particle theory (11). Similarly, there is a danger of an over-abundance of monopoles, a problem nowadays addressed by theories of cosmic inflation.

Cosmic strings, on the other hand, do not suffer from the same over-dominance problems as domain walls and monopoles. Over the years, Tom and co-workers including Turok (12), Lazarides and Shafi (13), Austin and Copeland (15), Hindmarsh (16), Vilenkin (17), Davis (21) and Vachaspati (25) considered models for structure formation seeded by cosmic strings. Tom’s vision has thereby provided an extraordinary link between the macroscopic and microscopic features of our Universe (see figure 6).
2.7. String solitons, M-theory and D-branes

The developments in Tom’s cosmic strings, non-perturbative macroscopic solutions of grand unified quantum field theories and perturbative superstrings, candidates for the fundamental microscopic constituents of matter incorporating gravity, proceeded along different lines—although in 1985 Edward Witten (ForMemRS 1999) speculated that fundamental strings produced in the early universe might be stretched to macroscopic scales, which, however, get diluted away by inflation (Witten 1985).

Things began to change with the realization that topological defects such as monopoles, strings and domain walls appear as solitonic solutions of the string equations, which, in the low energy limit, are those of supergravity. Loosely speaking, these solitons are non-singular field configurations that solve the source-free field equations, carry a non-vanishing topological ‘magnetic’ charge and are stabilized by a topological conservation law. They are intrinsically non-perturbative, with a mass depending on the inverse of the coupling constant. See the reviews (Duff et al. 1995, co-authored by one of us (MJD) while attending the six-month long workshop Topological defects organized by Tom at the Isaac Newton Institute, Cambridge, in 1994; Stelle 1997). They are to be compared and contrasted with the elementary field configurations, which are singular solutions of field equation with a source term and carry a non-vanishing Noether ‘electric’ charge.

Moreover, since string theory allows 10 spacetime dimensions, and its non-perturbative completion M-theory allows 11, they admit as solutions a menagerie of higher-dimensional objects in addition to the fundamental F-string, such as M-branes, Dirichlet-branes (D-branes) and Neveu–Schwarz-branes (NS-branes). They behave like strings if all but one
of their spatial dimensions are wrapped around the extra dimensions of spacetime. As Tom remarked (24): ‘string theory cosmologists have discovered cosmic strings lurking everywhere in the undergrowth.’ In particular Sarangi & Tye (2002) predicted the production of cosmic superstrings in the last stages of brane inflation, and Polchinski (2005) realized that the expanding Universe could not only stretch F- and D-strings to galactic size but, with large extra dimensions, they could be light enough to survive inflation. Tom continued to make significant contributions to this new cosmic stringy lease of life (24).

Cosmic strings are not the dominant source of density perturbations. Precision measurements of the cosmic microwave background (CMB) suggest that initial random Gaussian fluctuations, subsequently amplified by inflation, are the dominant source and that cosmic strings cannot account for more than about 10% of CMB structure. However, some such component could still have observable consequences. Gravitational lensing and gravitational radiation also provide possible observational features of cosmic strings.

2.8. Condensed matter systems

Wojciech Zurek pointed out that Tom’s description of topological defects also apply to the phase transition of normal fluid helium to superfluid helium (Zurek 1985). This became known as the ‘Kibble–Zurek mechanism’. Such a correspondence is found also in other condensed matter systems, for example, liquid crystals (20, 22). Indeed, Tom was involved in an experiment to search for defect formation in the B phase of He-3 where the experimental conditions mimic the cosmological Big Bang (18). His work inspired several experimental teams to search for defects formed in phase transitions in condensed matter systems; for example, Lancaster (He-3 and He-4), Grenoble and Helsinki (superfluid systems) and the US (liquid crystals). These effects have been experimentally confirmed in the context of vortex formation in superfluid helium 3 (18, 20, 22).

3. Leadership, mentoring and social responsibility

In 1970 Tom became professor of theoretical physics at Imperial, and held the position of head of the Department of Physics from 1983 to 1991. He was an outstanding teacher, and his undergraduate textbook Classical mechanics (now in a fifth edition co-authored with Frank H. Berkshire) is widely regarded as the best of its kind. He was also one of the winners of a competition set by the then Secretary of State William Waldegrave to explain the Higgs boson to the general public on a single sheet of A4.

Tom was concerned about the nuclear arms race, and took leading roles in several organizations promoting the social responsibility of science. These included the British Society for Social Responsibility in Science, Scientists against Nuclear Arms, Scientists for Global Responsibility and the Martin Ryle Trust.

4. Personality

Brilliant physicist, inspirational teacher and leader though he was, Tom conducted himself with an amazing degree of modesty and humility. He also showed unfailing kindness to the colleagues and students with whom he came into contact. Tom and his wife Anne frequently invited the whole Theory Group to parties at their home in Richmond.
A colleague of Tom at Imperial College, Anne-Christine Davis, now emeritus professor of mathematical physics at the Department of Applied Mathematics and Theoretical Physics (DAMTP), Cambridge, writes (personal communication 2020): ‘Tom was a mentor to many younger physicists and a strong supporter of women physicists, helping not just those in his immediate vicinity but others in the wider community. At times when one needed a sympathetic ear, or good advice, Tom was there and was at times instrumental in helping female physicists. His dedication to helping and mentoring younger physicists and in particular female physicists was recognized with the 2005 Mentoring lifetime achievement award, NESTA/Nature. Tom also had a good sense of humour, sometimes dry humour, sometimes even wicked! At Imperial Tom liked most to be surrounded by others in his research group, in particular the younger members and would make a point of joining them for lunch and tea rather than sitting with senior colleagues.’

In a moving eulogy, his son Robert (personal communication 2020) provided some anecdotes, giving us more glimpses of Tom’s personality. One recalls him being introduced to someone, who said: ‘I remember you teaching me 30 years ago, and I’ve just come back to finish my PhD.’ Tom congratulated him, shook his hand, and afterwards asked who that had been, to be told that it was Queen guitarist Brian May. Another recalls that while travelling on a bus in his 70s, a pickpocket stole his wallet. Tom leapt off the bus and chased him down the street, eventually forcing him to drop it.

Tom was a passionate supporter of many causes, frequently ending up running the organizations he joined (Scientists against Nuclear Arms and the Richmond Rambling Association, to name but two). A little known fact is that Tom was also an accomplished dancer.

5. Discovery of the Higgs boson: Nobel dilemma

As we have seen, one of Tom’s most important pieces of work was his study of the spontaneous symmetry-breaking mechanism whereby vector particles can acquire a mass accompanied by the appearance of a massive scalar boson. This mechanism, put forward independently by Brout & Englert (1964), by Higgs (1964), and by Guralnik, Hagen & Kibble (2) in 1964 lies at the heart of the Standard Model and all modern unified theories of fundamental particles (23). It was vindicated in 2012 by the discovery of the Higgs boson at CERN; the massive W and Z vector bosons having already been discovered at CERN in 1982 and 1983, respectively. Figure 7 shows Tom standing next to the Compact Muon Solenoid (CMS), which, together with ATLAS, detected the Higgs boson. This presented a much-debated dilemma for the Nobel Committee, notwithstanding Brout’s demise in 2011, as the Prize can be awarded to at most three recipients. However, Tom was also sole author of a 1967 paper, focusing on the non-abelian generalization (4). Indeed, in early papers, the mechanism was known as the ‘Higgs–Kibble effect’; see, for example, the paper by ’t Hooft and Veltman (’t Hooft & Veltman 1972) for which they won the 1999 Nobel Prize. Higgs and Kibble were the joint winners of the 1981 Hughes Medal of the Royal Society and the 1984 Rutherford Medal and Prize of the Institute of Physics.

Tom’s 1967 paper laid out the complete mathematical structure of spontaneous symmetry breaking, settling the issue of which gauge fields become massive and which remain massless. Private discussions with Salam informed his thinking on the problem. For a gauge field
theory model based on a non-abelian gauge group $G$ and a scalar Higgs sector producing spontaneous symmetry breaking down to a subgroup $H$, vector fields corresponding to the coset $G/H$ become massive, while the vectors corresponding to the unbroken subgroup $H$ remain massless. When the Standard Model came along shortly afterwards, Tom’s paper was seen to explain not only why the $W$ and $Z$ acquire a mass but, equally crucial, why the photon does not. His eightieth birthday was marked by Imperial College in March 2013 with a symposium day. In his evening public lecture, Nobel Laureate Steven Weinberg (Weinberg
2013) ended by saying: ‘Tom Kibble showed us why light is massless.’ Indeed Higgs said that Tom should have shared the 2013 Nobel Prize awarded to Englert and himself ‘because of what he wrote in 1967’ (Paull 2013). Kibble himself maintained a dignified modesty throughout in keeping with the honesty and integrity for which he was justly famous.

6. HONOURS AND AWARDS

1980 Fellow of the Royal Society
1981 Hughes Medal of the Royal Society (with P. W. Higgs)
1984 Rutherford Medal and Prize, Institute of Physics (with P. W. Higgs)
1991 Fellow of the Institute of Physics
1993 Guthrie Medal and Prize, Institute of Physics
1998 CBE
2005 Mentoring lifetime achievement award, NESTA/Nature
2009 Dirac Medal IOP
2010 Sakurai Prize for Theoretical Physics
2012 Royal Medal, Royal Society
2012 Honorary Fellow IOP
2012 Einstein Medal
2013 Knighthood
2014 Royal Medal, Royal Society of Edinburgh
2016 Isaac Newton Medal IOP

ACKNOWLEDGEMENTS

We thank Carlo Contaldi, Anne Davis, Graziela De Nadai, Philip Diamond, John Ellis, David Fairlie, Chris Isham, Robert Kibble, Alison Martin, Arttu Rajantie, Ray Rivers and Helen Wilson for help in preparing this memoir.

The portrait photograph was taken in 1980 by Godfrey Argent (© Godfrey Argent Studio).

AUTHOR PROFILES

Michael Duff

Michael Duff (FRS 2009) is emeritus professor of theoretical physics and senior research investigator at Imperial College London. He was formerly Abdus Salam professor of theoretical physics and principal of the Faculty of Physical Sciences. His time in the Theory Group overlapped with that of Tom Kibble for a total of 23 years. He is recipient of the 2004 Meeting Gold Medal, El Colegio Nacional, Mexico, the 2017 Paul Dirac Gold Medal and Prize, Institute of Physics, UK, and the 2018 Trotter Prize. He is a visitor to The Institute of Quantum Science and Engineering and was inducted as the Hagler Institute for Advanced Study 2019 Faculty Fellow, both at Texas A&M University.
Kellogg Stelle

Kelly Stelle is professor of theoretical physics at Imperial College London and was formerly head of the Theoretical Physics Group. He is the recipient of the 2006 Humboldt Research Award, from the Alexander von Humboldt Foundation, and the 2020 John William Strutt, Lord Rayleigh Medal and Prize, from the Institute of Physics. (Photograph courtesy of Imperial College London.)

REFERENCES TO OTHER AUTHORS

Anderson, P. W. 1963 Plasmons, gauge invariance and mass. *Phys. Rev.* 130, 439–442. (doi:10.1103/PhysRev.130.439)

Brout, R. & Englert, F. 1964 Broken symmetry and the mass of gauge vector mesons. *Phys. Rev. Lett.* 13, 321–323. (doi:10.1103/PhysRevLett.13.321)

Copeland, E., Myers, R. & Polchinski, J. 2004 Cosmic F- and D-strings. *J. High Energy Phys.* 6, 13. (doi:10.1088/1126-6708/2004/06/013)

Duff, M. J., Khuri, R. R. & Lu, J. X. 1995 String solitons. *Phys. Rep.* 259, 213–326. (doi:10.1016/0370-1573(95)00002-X)

Glashow, S. 1961 Partial symmetries of weak interactions. *Nucl. Phys.* 22, 579–588. (doi:10.1016/0029-5582(61)90469-2)

Goldstone, J., Salam, A. & Weinberg, S. 1962 Broken symmetries. *Phys. Rev.* 127, 965–970. (doi:10.1103/PhysRev.127.965)

Hanneshottir, H. & Schwartz, M. D. 2020 An S-matrix for massless particles. *Phys. Rev. D* 101, 105001. (doi:10.1103/PhysRevD.101.105001)

Higgs, P. 1964 Broken symmetries, massless particles and gauge fields. *Phys. Lett.* 12, 132–133. (doi:10.1101/0031-9163(64)91136-9)

Nambu, Y. 1960 Dynamical theory of elementary particles suggested by superconductivity. In *Proc. 1960 Annu. Int. Conf. High Energy Physics* (ed. E. C. G. Sudarshan, J. H. Tinlot & A. C. Melissinos), pp. 858–866. Rochester, NY: University of Rochester.

Nambu, Y. & Jona-Lasinio, G. 1961a Dynamical model of elementary particles based on an analogy with superconductivity: I. *Phys. Rev.* 122, 345–358. (doi:10.1103/PhysRev.122.345)

Nambu, Y. & Jona-Lasinio, G. 1961b Dynamical model of elementary particles based on an analogy with superconductivity: II. *Phys. Rev.* 124, 246–254. (doi:10.1103/PhysRev.124.246)

Paull, N. 2013 Peter Higgs: impact of God particle 'exaggerated'. *The Scotsman*, 13 November. (https://www.scotsman.com/news/uk-news/peter-higgs-impact-god-particle-exaggerated-1553389)

Polchinski, J. 2005 Cosmic superstrings revisited. *Int. J. Mod. Phys. A* 20, 3413–3415. (doi:10.1142/S0217751X05026886)

Polyakov, A. M. 1974 Particle spectrum in the quantum field theory. *JETP Lett.* 20(6), 194–195.

Salam, A. 1968 Weak and electromagnetic interactions. In *Elementary particle physics: relativistic groups and analyticity: Proc. 8th Nobel Symposium* (ed. N. Svartholm), p 367. Stockholm, Sweden: Almqvist and Wiksell.

Salam, A. & Ward, J. 1964 Electromagnetic and weak interactions. *Phys. Lett.* 13, 168–171. (doi:10.1016/0031-9163(64)90711-5)
Sarangi, S. & Tye, S. H. 2002 Cosmic string production towards the end of brane inflation. *Phys. Lett. B* **536**(3–4), 185–192. (doi:10.1016/S0370-2693(02)01824-5)

Shaw R. 1955 PhD thesis, Cambridge University.

Stelle, K. S. 1997 BPS branes in supergravity. In *High energy physics and cosmology: Proc. Summer School, Trieste, Italy, 2 June–4 July, 1997*, pp. 29–127. Report number: CERN-TH-98-80, IMPERIAL-TP-97-98-30. Geneva, Switzerland: CERN.

’t Hooft, G. & Veltman, M. 1972 Regularization and renormalization of gauge fields. *Nucl. Phys. B* **44**(1), 189–213. (doi:10.1016/0550-3213(72)90279-9)

’t Hooft, G. 1974 Magnetic monopoles in unified gauge theories. *Nucl. Phys. B* **79**, 276–284. (doi:10.1016/0550-3213(74)90486-6)

Weinberg, S. 1967 A model of leptons. *Phys. Rev. Lett.* **19**, 1264–1266. (doi:10.1103/PhysRevLett.19.1264)

Weinberg, S. 2013 Symposium evening lecture at Imperial College London. (http://www3.imperial.ac.uk/newsandevents/epsg/prggrp/imperialcollege/eventssummary/event_21-1-2013-16-35-46)

Witten, E. 1985 Cosmic superstrings. *Phys. Lett. B* **153**, 243–246. (doi:10.1016/0370-2693(85)90540-4)

Yang, C. N. & Mills, R. L. 1954 Conservation of isotopic spin and isotopic gauge invariance. *Phys. Rev.* **96**, 191–195. (doi:10.1103/PhysRev.96.191)

Zurek, W. 1985 Cosmological experiments in superfluid helium? *Nature* **317**, 505–508. (doi:10.1038/317505a0)

### Bibliography

The following publications are those referred to directly in the text. A full bibliography is available as electronic supplementary material at https://doi.org/10.6084/m9.figshare.c.5348674.

1. 1961 Lorentz invariance and the gravitational field. *J. Math. Phys.* **2**, 212–221. (doi:10.1063/1.1703702)

2. 1964 (With G. S. Guralnik & C. R. Hagen) Global conservation laws and massless particles. *Phys. Rev. Lett.* **13**, 585–587. (doi:10.1103/PhysRevLett.13.585)

3. (With L. S. Brown) Interaction of intense laser beams with electrons. *Phys. Rev.* **133**, A705–A719. (doi:10.1103/PhysRev.133.A)

4. 1965 Symmetry-breaking in non-Abelian gauge theories. *Phys. Rev.* **155**, 1554–1561. (doi:10.1103/PhysRev.155.1554)

5. 1968 Coherent soft-photon states and infrared divergences. I: Classical currents. *J. Math. Phys.* **9**, 315–324. (doi:10.1063/1.16645820)

6. Coherent soft-photon states and infrared divergences. II: Mass-shell singularities of Green’s functions. *Phys. Rev.* **173**, 1527–1535. (doi:10.1103/PhysRev.173.1527)

7. Coherent soft-photon states and infrared divergences. III: Asymptotic states and reduction formulas. *Phys. Rev.* **174**, 1882–1901. (doi:10.1103/PhysRev.174.1882)

8. Coherent soft-photon states and infrared divergences. IV: The scattering operator. *Phys. Rev.* **175**, 1624–1640. (doi:10.1103/PhysRev.175.1624)

9. 1976 The topology of cosmic domains and strings. *J. Phys. A Math. Gen.* **9**, 1387. (doi:10.1088/0305-4470/9/8/029)

10. 1980 Some implications of a cosmological phase transition. *Phys. Rep.* **67**, 183–199. (doi:10.1016/0370-1573(80)90091-5)

11. 1982 (With N. Turok) Self-intersection of cosmic strings. *Phys. Lett. B* **116**, 141–143. (doi:10.1016/0370-2693(82)90093-5)

12. (With G. Lazarides & Q. Shafi) Walls bounded by strings. *Phys. Rev. D* **26**, 435–439. (doi:10.1103/PhysRevD.26.435)

13. 1986 (With K. S. Stelle) Gauge theories of gravity and supergravity. In *Progress in quantum field theory. Festschrift for H. Umezawa* (ed. H. Ezawa & S. Kamefuchi), pp. 57–81. Amsterdam, The Netherlands: North-Holland.

14. 1993 (With D. Austin & E. Copeland) Evolution of cosmic string configurations. *Phys. Rev. D* **48**, 5594–5627. (doi:10.1103/PhysRevD.48.5594)
(15) 1995 (With M. B. Hindmarsh) Cosmic strings. *Rep. Progr. Phys.* **58**, 477–562. (doi:10.1088/0034-4885/58/5/001)

(16) 1995 (With A. Vilenkin) Phase equilibration in bubble collisions. *Phys. Rev. D* **52**, 679–688. (doi:10.1103/PhysRevD.52.679)

(17) 1996 (With V. M. H. Ruutu, V. B. Eltsov, A. J. Gill, M. Krusius, Y. G. Makhlin, B. Plaçais, G. E. Volovik & W. Xu) Vortex formation in neutron-irradiated superfluid 3He as an analogue of cosmological defect formation. *Nature* **382**, 334–336. (doi:10.1038/382334a0)

(18) 1999 Recollections of Abdus Salam at Imperial College. In *The Abdus Salam Memorial Meeting: Trieste (Italy), 19–22 November 1997* (ed. J. Ellis, F. Hussain, G. Thompson, M. Virasoro), pp. 1–11. Singapore: World Scientific.

(19) 2003 Symmetry breaking and defects. In *Patterns of symmetry breaking. Proc. NATO Advanced Study Institute, Cracow, September 2002* (ed. H. Arodz, J. Dziarmaga & W. H. Zurek), pp. 3–36. Dordrecht, The Netherlands: Springer. (doi:10.1007/978-94-007-1029-0)

(20) 2005 (With A.-C. Davis) Fundamental cosmic strings. *Contemp. Phys.* **46**, 313–322. (doi:10.1080.00107510500165204)

(21) 2007 Phase transition dynamics in the lab and the universe. *Phys. Today* **60** (9), 47–53. (doi:10.1063/1.2784684)

(22) 2009 Englert–Brout–Higgs–Guralnik–Hagen–Kibble mechanism. *Scholarpedia* **4** (1), 644.

(23) 2010 (With E. J. Copeland) Cosmic strings and superstrings. *Proc. R. Soc. A* **466**, 623–657. (doi:10.1098/rspa.2009.0591)

(24) 2015 (With T. Vachaspati) Monopoles and strings. *J. Phys.* **G42**, 094002. (doi:10.1088/0954-3899/42/9/094002)