Twist grain boundary (TGB) states of chiral liquid crystalline bent-core mesogens†

Hale Ocakı,ab Belkiz Bilgin-Eran,b Dilek Güzeller,d Marko Prehmab and Carsten Tschierskea,*

4-Cyanoresorcinol derived bent-core molecules with a chiral (S)-2-methylbutoxy chain form liquid crystalline phases with TGBA- and TGBC-like structures at the transition from cybotactic nematic via SmA to SmC phases.

Chirality has huge effects on molecular self-assembly.1 Soft matter, especially liquid crystalline (LC) phases,2,3 can easily be affected by chirality leading to helical superstructures.4 Molecules with high helical twisting power can even give rise to the frustration of the fundamental structures of LC self-assembly, thus providing new superstructures with a higher level of complexity3 and new emerging properties. Examples are the three-dimensional lattices of defects in the blue phases which are of importance as photonic band gap materials and represent candidates for new generations of extremely fast switching electro-optical devices,5 twist grain boundary (TGB) phases,6 representing analogues of the Abrikosov flux phase of type-II superconductors7 and several other complex structures.5,8

TGB phases, formed by chiral rod-like mesogens, represent helical superstructures with a helix axis parallel to the layer planes thus disrupting the layers to give blocks separated by screw dislocations (Fig. 1). TGB phases can be divided into TGBA phases with an orthogonal organization of the molecules with respect to the planes of these blocks and several different modes of TGBC phases with a tilted organization of the molecules.6 In the recent two decades bent-core (BC) molecules9 have attracted special attention as new LC materials. This is mainly due to their capability to form LC phases with polar order9 and the spontaneous formation of chiral superstructures though the molecules are achiral.9–12 An interesting point concerns the interaction of these chiral superstructures with the molecular chirality provided by stereogenic units. For example, enhancement of chirality was observed by doping chiral LC phases with achiral bent-core mesogens, which lead to the induction of highly frustrated blue phases with helical organization in all three directions instead of only one in the cholesteric phases.5 It was proposed that “highly chiral” conformations11,13 of the bent aromatic cores lead to a high twisting power.† For the same reason broad regions of blue phases were recently found for optically active BCLCs.14

Herein we report the first observation of TGB-like structures for the LC phases of the BC mesogens (S)-1/n (n = 12, 14) involving a 4-cyanoresorcinol derived BC,14b,15 combined with a chiral (S)-2-methylbutoxy group. In order to identify the specific chirality effects, one of the compounds was also synthesized in racemic form (rac-1/12). The synthetic strategy of these BC compounds is outlined in Scheme 1 and the experimental procedures are described in the ESL.†

The observed LC phases, their transition temperatures and associated enthalpies are collated in Table 1. Upon cooling the racemic mixture rac-1/12 from the isotropic liquid (Iso) a nematic (N) phase is formed at first which is then replaced by a lamellar phase (SmA) at $T \leq 96$ °C. In the SmA phase there is on average a non-tilted organization of the molecules in layers, whereas at the transition to the SmC phase at $T = 88$ °C a uniform tilt develops. In calorimetric investigations besides the dominating melting peak

† Electronic supplementary information (ESI) available: Synthesis of compounds, analytical data, DSC traces, additional XRD patterns, textures and other data. See DOI: 10.1039/c5cc01592h
It appears that the LC phases can also be easily investigated (see Table 1).

The typical textures observed for the LC phases of (S)-1/12 are the same as those found for rac-1/12. For (S)-1/14 a very similar behaviour to those for (S)-1/12 is observed, only the transitions between the LC phases occur at a bit higher temperatures (Table 1). The typical textures observed for the LC phases of (S)-1/n (n = 12, 14) are shown in Fig. 2. A cholesteric oily-streak texture indicates the presence of a helical superstructure in the chiral nematic phases (N*) (Fig. 2a).5 The appearance of the LC phases occurring below N* strongly depends on the alignment conditions. Homeotropic anchoring (layers parallel to the surfaces of the glass substrates) leads to filament textures, as characteristic for TGBA phases16 (Fig. 2b). The filaments rapidly disappear and the textures become uniformly dark, indicating the transition to an optically uniaxial SmA phase. Upon further cooling, a gray and low birefringent texture develops in these homeotropic samples, indicating a transition to a SmC* phase, occurring for (S)-1/12 at about the same temperature as observed for the SmA-SmC transition of the racemate.§ In a cell with planar surface anchoring (layers perpendicular to the substrate surfaces; polyimide-coated ITO cell, 10 μm) the cholesteric oily streak texture changes at the N*-SmA transition to a planar texture composed of differently coloured areas,
corresponding to different twist states (Fig. 2c). This texture, which is typical for TGBA phases, is retained upon further cooling (see Fig. S2, ESI†) and at the SmA/SmC phase transition temperature the formation of a grid-like pattern is observed (Fig. 2d), similar to textures of TGBC phases.4,6,17 The texture continuously changes upon further cooling (see Fig. S3, ESI†). Therefore it is not clear if the TGBC structure is retained or a slow transition to a SmC-phase takes place upon further cooling (therefore it is denoted as SmC*/TGBC in Table 1).

XRD investigations were performed for (S)-1/12 and rac-1/12 in thin capillaries at a magnetic field B ∼ 1 T after slow cooling (0.1 K min⁻¹) from the isotropic liquid. In the whole investigated temperature range between 60 and 110 °C the wide angle scattering is diffuse indicating the presence of LC phases without long range order between individual molecules (Fig. 3a and b). In the small angle range there is a diffuse scattering in the N and N* phases, exceeding the wide angle scattering and thus confirming cybotactic nematic phases composed of smectic clusters.15 For rac-1/12 the small angle scattering maximum is on the meridian, indicating that the molecules have no uniform tilt in these domains, i.e. the nematic phases are composed of SmA clusters (N_{cyt}), which is rarely found for BC molecules.15 The intensity of the small angle scattering increases further at the transition to the SmA phase (Fig. 3d). The d-spacings of 4.05–4.10 nm, corresponding to ~0.85 molecular length (L_{mol} = 4.8 nm in the most stretched conformation, see Fig. S8, ESI†), are in line with a single layer structure of the cybotactic clusters and layers. The d = f(T) plot for rac-1/12 indicates an increase of d in the N and SmA phases due to an increasing packing density, leading to alkyl chain stretching. A slight decrease of d starts at around the SmA-SmC transition, in line with the onset of uniform tilt at this transition (Fig. 3c). The tilt in the SmC phase is very small (<5°) as estimated for rac-1/12 from the angle between the diffuse wide angle scattering maxima on the equator and the layer reflections on the meridian (XRD tilt). The optical tilt is about 10–15° as estimated from the planar textures of rac-12 (see Fig. S1g and h, ESI†), indicating that the main contribution comes from the tilt of the aromatic cores.

The LC phases of (S)-1/12 (Fig. 3b) have approximately the same d-values as those of rac-1/12. However, in the N° phase the small- and wide-angle scatterings form closed rings, indicating the absence of uniform alignment due to the presence of the helical superstructure. At the transition to TGBA the rings condense to crescent like maxima on the meridian (small angle) and on the equator (wide angle), respectively, in line with the formation of layers. There is a broader angular χ-distribution of the scattering in the smectic phases of (S)-1/12 compared to the sharper peak of rac-1/12 (Fig. 3a and b). However, there is no distribution of the scatterings on a closed ring as it would be expected for TGB phases. So, there seems to be a strong effect of the conditions on the actually observed phase structure. As shown above, planar anchoring stabilizes the TGB helix, whereas homeotropic anchoring has the opposite effect; it stabilizes the layers and tends to remove the TGB helix. Hence, the TGB structures are not considered as “phases”, but as helically deformed states of the underlying smectic phases. The reduced influence of surface anchoring in the capillaries used for XRD, and the alignment of the molecules by the magnetic field, might suppress TGB-helix formation. The broader angular distribution of the scattering of (S)-1/12 might be due to the remaining helical layer distortion or to the less efficient alignment obtained by cooling from the N° phase.

There is no polarization current response under a triangular wave field (PI coated ITO cells 10 μm) in the temperature range of all mesophases for the racemic mixtures as well as for the (S)-enantiomers up to the maximum available value of ±28 V μm⁻¹, indicating the absence of polar order.¶

As an additional interesting point, it should be noted that chiral domains are visible in homeotropic alignment in the range of the SmC phase of the racemic compound rac-1/12. These domains can be identified by slight rotation of the analyzer either clockwise or anti-clockwise out of the precise 90° orientation, which reverses the brightness of these domains (Fig. 4a–c). Rotating the sample between the polarizers does not lead to such a change in brightness, thus excluding tilt director alignment as the origin of this effect (Fig. S6, ESI†). In addition, though domains can also be recognized in the homeotropic SmC* phase of (S)-1/12, in this case there is no visible effect of the orientation of the polarizer (Fig. 4d–f). As there is no polar order in the SmC phase the origin of chirality must be due to the segregation of chiral molecular conformers,13 which is additionally supported by surface interaction.|| This indicates a strong chirality of the molecular conformations and a strong
coupling between them. The cybotactic nature of the nematic phase and the absence of clear transition enthalpies at the phase transitions indicate a nearly continuous growth of the coherence length of the smectic clusters throughout the $N_{cyb}$-SmA-SmC transitions (Fig. 3d, Table S1, ESI†). This indicates that already in the nematic phase the cybotactic clusters are relatively large and continue to increase in the smectic phases. Thus, in the SmA and SmC phases there is a remaining layer distortion providing soft layers which can be easily deformed into a TGB-like superstructure (Fig. 1). Therefore, the weak helical twisting power of the (S)-2-methylbutyl group can provide a sufficiently strong chirality effect such that TGB states are formed at the N*-SmA transition. These can be further stabilized over broader temperature ranges by planar surface anchoring.

In summary, the first observation of TGB structures in meso-phases of bent-core mesogens is reported.** These are formed by molecules with a (S)-2-methylbutoxy stereogenic centre, known to have only weak helical twisting power. However, strong chirality of the molecular conformers, an imperfect layer structure and surface anchoring can stabilize the TGB states.

The work was supported by DFG (Ts 39/24-1); H. O. is grateful to the Alexander von Humboldt Foundation for a research fellowship at Martin Luther University, Halle-Wittenberg; B. B.-E. is grateful to the Alexander von Humboldt Foundation for financial support toward LC research.

Notes and references

† That no switching peak could be observed even in the SmC* phases could be due to the known weak effect of the (S)-2-methylbutyl group on polar order, leading to very small values of spontaneous polarization.1,14,16 Also optically, no switching can be observed; this confirms that there is no switching or it takes place by rotating around the molecular long axis.

¶ This kind of chiral SmC phase was previously reported for 4-cyanoresorcinol based BC molecules with azobenzene wings, and a local SmC$_P$ structure was assumed to be responsible for chirality.1,6,19 However, the recent observation of chiral segregation in isotropic liquids3 supports the possibility of spontaneous chiral segregation in SmC phases with a dense packing of the aromatic cores. This chiral segregation provides a helical distortion of the layers, thus giving rise to layer distortion and leading to amplification of small chirality effects.

** TGB-like twist states and TGB phases were observed for mixtures of chiral rod-like molecules with achiral bent-core molecules19 or hockey-stick compounds20 and for chiral dimesogens with odd spacers.1 A TGB-like structure has also been discussed as a possible organization in the dark conglomerate (DC) phases of achiral BC molecules21 but only recently a DC-like random grain boundary phase of achiral hockey-stick LC was reported.21b

1 D. B. Amabilino, *Chirality at the Nanoscale*, Wiley-VCH, Weinheim, 2009.
2 *Handbook of Liquid Crystals*, ed. J. W. Goodby, P. J. Collings, T. Kato, C. Tscharke, H. F. Gleeson and P. Raynes, Wiley-VCH, Weinheim, 2014.
3 C. Tscharke, *Angew. Chem., Int. Ed.*, 2013, 52, 8828–8878.
4 I. Dierking, *Symmetry*, 2014, 6, 444–472.
5 *Chirality in liquid crystals*, ed. H.-S. Kietzerow and C. Bahr, Springer, New York, NY, 2001.
6 (a) J. W. Goodby, M. A. Auwa, S. M. Stein, E. Chin, R. Pink and J. S. Patel, *Nature*, 1989, 337, 449–452; (b) J. W. Goodby, *Curr. Opin. Colloid Interface Sci.*, 2002, 7, 326–332; (c) M. Brunet, L. Navailles and N. A. Clark, *Eur. Phys. J. E: Soft Matter Biol. Phys.*, 2002, 7, 5–11.
7 S. R. Renn and T. C. Lubensky, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1988, 38, 2132–2147.
8 H. T. Nguyen, M. Ismaili, N. Isaert and M. F. Achard, *J. Mater. Chem.*, 2004, 14, 1560–1566.
9 (a) R. A. Reddy and C. Tscharke, *J. Mater. Chem.*, 2006, 16, 907–961; (b) H. Takezoe and Y. Takashina, *Ipn. J. Appl. Phys.*, 2006, 45, 597–625; (c) A. Eremin and A. Jakli, *Soft Matter*, 2013, 9, 615.
10 L. E. Hough, M. Spannuth, M. Nakata, D. A. Coleman, C. D. Jones, G. Dantlgraber, C. Tscharke, J. Watanabe, E. Körblová, D. M. Walba, J. E. MacManne, M. A. Glaser and N. A. Clark, *Science*, 2009, 325, 452.
11 H. Takezoe, *Top. Curr. Chem.*, 2012, 318, 303–330.
12 V. P. Panov, R. Balachandran, J. K. Vij, M. G. Tamba, A. Kohmeier and G. H. Mehl, *Appl. Phys. Lett.*, 2012, 101, 234106.
13 C. Dressel, T. Reppe, M. Prehm, M. Brautzsch and C. Tscharke, *Nat. Chem.*, 2014, 6, 971–977.
14 (a) H.-C. Jeong, S. Aya, S. Kang, F. Araoka, K. Ishikawa and H. Takezoe, *Liq. Cryst.*, 2013, 40, 951–958; (b) H. Oçak, B. Bilginen, M. Prehm, S. Schymura, J. P. F. Lagerwall and C. Tscharke, *Soft Matter*, 2011, 7, 8266–8280.
15 (a) L. Kovalenko, M. W. Schröder, R. A. Reddy, S. Diele, G. Pelzl and W. Weissflog, *Liq. Cryst.*, 2005, 32, 857–865; (b) C. Keith, A. Lehmann, U. Baumeister, M. Prehm and C. Tscharke, *Soft Matter*, 2010, 6, 1704–1721.
16 I. Dierking and S. T. Lagerwall, *Liq. Cryst.*, 1999, 26, 83–95.
17 C. V. Yelamagadda, A. S. Achalkumar, N. L. Bonde and A. K. Prajapati, *Chem. Mater.*, 2006, 18, 1076–1078.
18 M. Allassar, M. Prehm, M. Nagaraj, J. K. Vij and C. Tscharke, *Adv. Mater.*, 2013, 25, 2186–2191.
19 P. Archer and I. Dierking, *Liq. Cryst.*, 2006, 33, 257–265.
20 V. Novotna, M. Glogarova, V. Kozmik, J. Svoboda, V. Hamplova, M. Kaspar and D. Pociecha, *Soft Matter*, 2013, 9, 647–653.
21 (a) S. K. Lee, L. Shi, M. Tokita and J. Watanabe, *J. Phys. Chem. B*, 2005, 112, 6762–6766; (b) D. Chen, H. Wang, M. Li, M. A. Glaser, J. E. MacLennan and N. A. Clark, *Soft Matter*, 2014, 10, 9105–9109.