AN EXPANDED VERY LARGE ARRAY SEARCH FOR WATER MEGAMASER EMISSION IN THE SUBMM GALAXY SMM J16359+6612 AT z = 2.5

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ABSTRACT

Using the Expanded Very Large Array, we have conducted a search for 22.2 GHz H2O megamaser emission in the strongly lensed submm galaxy, SMM J16359+6612 at z = 2.517. This object is lensed into three components, and after a correction for magnification emission is applied to its submm-wavelength flux density, it is typical of the bulk of the high-redshift, submm galaxy population radio flux densities for the 850 μm extragalactic background (S850,μm ~ 1 mJy). We do not detect any H2O megamaser emission, but the lensing allows us to place an interesting constraint on the luminosity of any megamaser present, L850 < 5030 Lsun for an assumed linewidth of 80 km s⁻¹. Because the far-infrared luminosity in submm galaxies is mainly powered by star formation, and very luminous H2O megamasers are more commonly associated with quasar activity, it could be that blind searches for H2O megamasers will not be an effective means of determining redshifts for less-luminous members of the submm galaxy population.

Key words: cosmology: observations – galaxies: starburst – gravitational lensing – masers

1. INTRODUCTION

Single-dish submm/mm-wavelength surveys have revealed a population of massive, dust-enriched, far-infrared (FIR) luminous star-forming galaxies in the early Universe (Smail et al. 1997; Hughes et al. 1998; Barger et al. 1998; Bertoldi et al. 2000). Since their discovery, redshift determination for individual members of this submm galaxy (SMG) population has been hindered by the lack of sufficiently accurate astrometry needed to identify multiwavelength counterparts, limited by the coarse angular resolution of single-dish submm/mm telescopes (∼10⁶−20′′) and the faintness of SMGs at optical wavelengths. For the brightest members of the SMG population, deep radio interferometry imaging has provided a powerful solution to the former problem of identifying counterparts (e.g., Ivison et al. 1998, 2000, 2002, 2007; Smail et al. 2000; Webb et al. 2003; Clements et al. 2004; Borys et al. 2004; Dannerbauer et al. 2004), resulting in the ability to measure spectroscopic redshifts for the proposed optical counterparts to the members of the SMG population. Such surveys find a median spectroscopic redshift for bright SMGs (S850,μm > 5 mJy) of z ∼ 2.2 (Chapman et al. 2003, 2005). However, between 65% and 80% of bright SMGs have secure radio counterparts (e.g., Ivison et al. 2005, 2007), so a significant fraction may lie at z > 3.5, given the radio-wavelength selection effects. The uncertainty in the redshift distribution of the fainter SMG population (< 5 mJy at 850 μm) makes it difficult to complete the luminosity function for SMGs and, hence, difficult to accurately constrain the obscured star-formation rate of the Universe at the highest redshifts.

Alternative methods of redshift determination for SMGs have been explored, including photometric redshift techniques, which rely on the measured continuum flux densities at submm-to-radio wavelengths (Hughes et al. 1998; Carilli & Yun 1999, 2000; Aretxaga et al. 2003, 2005). These techniques have proven to be accurate to ±0.3 (rms) for SMGs at z ≤ 3.5, when three or more FIR-to-radio flux densities are measured (Aretxaga et al. 2005). For those SMGs with well-constrained spectral energy distributions (SEDs), the predictions of such techniques have helped to guide the choice of tunings for broadband cm-wavelength searches for redshifted CO line emission (Wagg et al. 2007), which should be an effective means of measuring SMG redshifts given the intensity of molecular CO line emission in these objects (e.g., Neri et al. 2003; Greve et al. 2005). The first cm-wavelength search for redshifted low-J CO line emission (J = 2 − 1/1 − 0) in a SMG with no optical spectroscopic redshift was unsuccessful (Wagg et al. 2007), while a great deal of uncertainty remains in the expected intensity in the low-J CO lines. Given the expected high densities of the molecular gas in the interstellar media of these SMGs (Weiβ et al. 2005), it is possible that mm-wavelength searches for high-J CO line emission will yield more positive results. Such observations will soon be possible with the ultrawide bandwidth “Redshift Receiver” on the Large Millimeter Telescope (LMT; Pérez-Grovas et al. 2006).

Broadband spectral line observations of SMGs at lower frequencies (νobs < 10 GHz) may yield redshifts through the detection of OH or H2O maser emission (Townsend et al. 2001; Ivison 2006). Such observations have an advantage over higher-frequency CO line searches in that the bandwidth needed to cover the same redshift interval decreases with decreasing frequency. This means that the broadband capabilities of the National Radio Astronomy Observatory (NRAO) Green Bank Telescope (GBT®), or the Expanded Very Large Array (EVLA), could potentially be exploited to measure redshifts for SMGs via OH or H2O maser emission lines. The luminosity in the 22,235 GHz H2O maser line has been found to weakly depend

5 It has been argued that this same level of uncertainty can be achieved if one assumes the median redshift of the bright SMG sample for individual objects (Chapman et al. 2005).

6 The NRAO is operated by Associated Universities Inc., under a cooperative agreement with the National Science Foundation.
on the FIR luminosity in nearby galaxies (Castangia et al. 2008), while a stronger dependence of the OH maser line luminosity on the FIR luminosity has been demonstrated (Darling & Giovanelli 2002). However, given that the OH maser lines emit at rest frequencies of 1665.4018 and 1667.3590 MHz, for objects at \( z > 2 \), these lines are redshifted to the \( v_{\text{obs}} \lesssim 600 \text{ MHz} \) band, where the presence of radio frequency interference (RFI) for many existing facilities may make it difficult to identify the OH lines in objects whose redshifts are unknown. For these same redshifts, the H_2O maser line is redshifted to frequencies in the 5–8 GHz range, where RFI is less of a problem. It should then be more effective to conduct blind searches for redshifted H_2O maser emission when attempting to measure redshifts for SMGs. Before such a technique can be applied, it must first be demonstrated that H_2O masers exist in SMGs with known redshifts. To date, the highest redshift detection of H_2O maser emission is that of the OH maser line in a type 2 quasar at \( z = 0.66 \) (Barvains & Antonucci 2005). Ivison (2006) conducted a search for both OH and H_2O megamaser emissions in the strongly lensed, ultraluminous infrared quasar, APM 08279+5255 at \( z = 3.91 \); however, no line emission is detected.

Here, we conduct a search for H_2O maser emission in the lensed SMG, SMM J16359+6612 at \( z = 2.517 \) (Kneib et al. 2004, 2005). Given the large total magnification factor inferred by the gravitational lensing model (\( m \approx 45 \)), this object has an intrinsic flux density typical of the bulk of the SMG population responsible for the 850 \( \mu \text{m} \) background radiation (\( S_{850,\mu \text{m}} \approx 1 \text{ mJy} \)). This object is strongly gravitationally lensed, and is the only SMG to have been detected in HCN \( J = 1–0 \) line emission (Gao et al. 2007), implying the presence of dense gas. SMM J16359+6612 presents a good candidate in which to conduct a search for H_2O maser emission, and one of only two strongly lensed SMGs for which such a search is possible. Our overall aim is to determine if searches for this line can be used to obtain blind redshifts for SMGs in future surveys and for this pilot study, we chose to focus on less-luminous SMGs typical of the submm extragalactic background. Throughout this work, we assume \( H_0 = 71 \text{ km s}^{-1}, \Omega_M = 0.27, \) and \( \Omega_\Lambda = 0.73 \) (Spergel et al. 2007).

2. OBSERVATIONS AND DATA REDUCTION

We observed SMM J16359+6612 with the Very Large Array (VLA) on 2008 June 4 in its compact DnC configuration. The details of these observations are given in Table 1. At the source redshift of \( z = 2.517 \) (determined from multiple CO emission lines; Weiß et al. 2005), the 22.23,508 GHz H_2O maser line is redshifted to 6,322 GHz, which is now accessible with the new C-band receivers available through the EVLA project. The observations were configured to have a 391 kHz channel width, equivalent to 18.5 km s\(^{-1}\) velocity resolution over a bandwidth of 12.5 MHz (591 km s\(^{-1}\)). In total, \( \sim 9 \) hr of on-source integration time was obtained with the 15 EVLA antennas available, leading to a synthesized beamsize of \( \sim 10^\prime \). Due to the current C-band receiver temperatures at nonstandard frequencies, and a source of unknown noise associated with using 12.5 MHz of bandwidth with the VLA correlator, our final sensitivities are \( \sim 5 \times \) worse than what will be achievable using the full array of EVLA antennas with the new Wideband Interferometric Digital ARchitecture (WIDAR) correlator.

Data reduction was performed using the Astronomical Image Processing System (AIPS) of NRAO. Some “flagging” was required, which resulted in two of the 15 EVLA antennas being removed entirely from the analysis. Flux, bandpass, and phase calibrations were derived from the observations of 3C286 and 1642+689. After applying the calibration solutions to the target data, uniform-weighted dirty Stokes I cubes were produced.

3. RESULTS

The submm/mm emission from SMM J16359+6612 is lensed into three components, where the observed flux densities of components A, B, and C have been amplified by factors of 14, 22, and 9, respectively (Kneib et al. 2004). No 6.3 GHz continuum emission is detected from any of the three components in the continuum image above an apparent rms of \( \sim 117 \mu \text{Jy} / \text{beam} \). These limits on the 6.3 GHz continuum emission are consistent with the value of 175 \( \mu \text{Jy} / \text{beam} \) for the peak 6.3 GHz flux density of the brightest component, predicted from the 1.4 and 8.2 GHz detections of SMM J16359+6612 with the VLA and the Westerbork Synthesis Radio Telescope (WSRT) by Garrett et al. (2005).

Spectra were extracted at the positions of each of the lensed components determined from the centroids of the high-J CO line emission. Each spectrum was corrected for lensing amplification. From the three spectra, an

| Parameter                      | Value                  |
|--------------------------------|------------------------|
| Frequency                      | 6.322 GHz              |
| Pointing center (J2000)        | 16\(^{\circ}\)35\(^{\prime}\)44\(\alpha\), 46\(^{\circ}\)46\(\alpha\)24\(\alpha\) |
| Synthesized beam               | 10\(\alpha\)72 \times 9\(\alpha\)31, P.A. = -7\(\alpha\)4 |
| Channel spacing                | 391 kHz               |
| Line rms (391 kHz channel)     | 520 \(\mu\)Jy/beam     |
| Continuum rms                  | 117 \(\mu\)Jy/beam     |
average spectrum was created using the “demagnified” noise of each in the weighting (Figure 2). Although our channel rms is \( \sim 520 \, \mu \text{Jy/beam} \), by first demagnifying and then averaging the spectra extracted at the positions of the three components, we are able to reach a sensitivity equivalent to \( \sim 16 \, \mu \text{Jy/beam} \) per 391 kHz channel.

Inspection of the individual spectral line data cubes did not reveal any significant detection of H\(_2\)O megamaser emission in SMM J16359+6612, nor is any emission detected in the average spectra of the three components (Figure 2). Although we do not detect redshifted H\(_2\)O maser emission in SMM J16359+6612, our sensitivity is sufficient to derive strong constraints on the luminosity of any H\(_2\)O megamaser emission present. Typical linewidths of H\(_2\)O maser emission lines in FIR bright galaxies are in the range of 1–130 km s\(^{-1}\), with one example of a \( \sim 260 \, \text{km s}^{-1} \) line in Arp299 (Henkel et al. 2005). As such, we assume intermediate linewidths of 20, 40, 60, and 80 \( \text{km s}^{-1} \), in order to calculate our H\(_2\)O line luminosity limits following the definition given in Solomon & Vanden Bout (2005). Before correcting the noise of an individual component for lensing magnification, the \( \sigma \) line luminosity limits are \( L_{\text{H}_2\text{O}} < (8.7, 12.2, 15.0, \text{and} 17.3) \times 10^4 \, L_\odot \), for the assumed linewidths. Much stronger constraints are obtained when we correct for lensing and calculate these luminosity limits for the average spectrum: \( L_{\text{H}_2\text{O}} < 2652, 3751, 4594, \text{and} 5305 \, L_\odot \).

4. DISCUSSION

We can estimate whether our limits to the H\(_2\)O megamaser line luminosity in SMM J16359+6612 are consistent with the value expected from the FIR luminosity and the \( L_{\text{FIR}}-L_{\text{H}_2\text{O}} \) luminosity correlation suggested by Castangia et al. (2008) for nearby galaxies. We note that this correlation is weak, and a wide range of H\(_2\)O line luminosities appear to be valid. The unlensed 850 \( \mu \text{m} \) flux density of each image of SMM J16359+6612 is \( \sim 0.8 \, \text{mJy} \) (Kneib et al. 2004), which can be converted to a FIR luminosity, \( L_{\text{FIR}} \sim 1.5 \times 10^{12} \, L_\odot \), for an assumed dust temperature, \( T_d = 40 \, \text{K} \). From the \( L_{\text{FIR}}-L_{\text{H}_2\text{O}} \) luminosity relation for Galactic star-forming regions (Genzel & Downes 1979; Jaffe et al. 1981), we should expect an H\(_2\)O megamaser line luminosity, \( L_{\text{H}_2\text{O}} \sim 1500 \, L_\odot \). This prediction is broadly consistent with the limits we obtain here for the narrowest linewidths assumed. For comparison, the luminosity in the H\(_2\)O megamaser line detected in a type 2 quasar at \( z = 0.66 \) by Barvainis & Antonucci (2005) is \( 2.3 \times 10^4 \, L_\odot \).

We propose possible explanations for our nondetection of H\(_2\)O megamaser emission in SMM J16359+6612. Although it is unlikely, it could be that the limited velocity covered by our choice of bandwidth (\( \sim 574 \, \text{km s}^{-1} \)) is too small, as the maser emission maybe at a significantly different redshift from that of the CO emission lines (e.g., Wilner et al. 1999). However, the more probable reason for our nondetection is that the molecular medium within SMM J16359+6612 is not conducive to the formation of megamaser emission. Neufeld et al. (1994) argued that luminous water maser emission can be excited by X-ray photons from an active galactic nucleus (AGN). Strong H\(_2\)O maser emission would arise in gas with temperatures of \( \sim 400–1000 \, \text{K} \) and pressures of \( P/k = 10^{11} \, \text{cm}^{-3} \, \text{K} \), irradiated by intense X-ray emission. Such AGN-powered circumnuclear H\(_2\)O megamasers can attain much larger line luminosities than the interstellar masers found in starburst galaxies, which are consequently termed “kilomasers.” X-ray studies of SMGs confirm that while \( \sim 28–50 \% \) do harbor an AGN, such AGN activity may only contribute to a small fraction of their enormous FIR luminosities (e.g., Alexander et al. 2005). Given the rarity, and lower luminosities of the AGN in SMGs, it is likely that the dense molecular gas in SMM J16359+6612 inferred from the HCN line emission is most likely associated with high-mass star formation. Extragalactic regions of star formation may give rise to water maser emission; however, this emission is generally weaker than that observed in galaxies with a luminous AGN. Similar is the case with the nearby galaxy, NGC 520, whose kilomaser line luminosity is lower than that expected from its FIR luminosity (Castangia et al. 2008). A third possibility for our nondetection of H\(_2\)O megamaser emission is that the emission is variable. Argon et al. (2007) found that the H\(_2\)O maser emission in the nearby galaxy, NGC 4258, varied by a factor of \( \sim 12 \) over a 3 year period.

This nondetection of H\(_2\)O megamaser emission in SMM J16359+6612 has consequences for the use of this line as a redshift indicator in future studies of less-luminous SMGs typical of the submm extragalactic background. Given the strong gravitational lensing amplification of this object, making it one of the brightest of the currently detected SMGs, J16359 should be one of the most likely SMGs to exhibit detectable H\(_2\)O megamaser emission. It has recently been suggested that the H\(_2\)O maser detection rate is weakly correlated with a galaxy’s FIR luminosity (Castangia et al. 2008). If this is the case, then it maybe that similarly bright, unlensed SMGs maybe more likely to contain H\(_2\)O megamasers. However, mounting evidence suggests that the brightest (\( S_{850 \mu \text{m}} \gtrsim 15 \, \text{mJy} \)) and, subsequently, most luminous SMGs maybe at the highest redshifts (e.g., Younger et al. 2007; Wang et al. 2007), and so cosmological dimming of the line intensity would necessitate large integration times (more than 100 hr), even with the full EVLA. Such searches would be best carried out with the future Square Kilometer Array (SKA).

5. SUMMARY

We have conducted the first search for 22 GHz water megamaser emission in a high-redshift SMG, by observing the strongly lensed, SMM J16359+6612 at \( z = 2.517 \), with the currently available EVLA antennas. No emission is detected, resulting in a 3\( \sigma \) luminosity limit of \( < 5305 \, L_\odot \) for an assumed linewidth of 80 km s\(^{-1}\). With an optimally working, completed
EVLA, we would have achieved $\sim 5 \times$ higher sensitivity, reaching down to the kilomaser luminosity range. Our nondetection suggests that luminous H$_2$O megamasers may not be present to provide reliable redshift estimates for the less-luminous members of the SMG population; however, future surveys of lensed SMGs with secure CO line redshifts are needed to strengthen this conclusion.

We thank the NRAO staff involved in the EVLA project for making these observations possible. J.W. and C.C. are grateful for support from the Max-Planck Society and the Alexander von Humboldt Foundation. We thank the anonymous referee for helpful suggestions on the original manuscript.

REFERENCES

Alexander, D. M., et al. 2005, Nature, 434, 738
Aretxaga, I., Hughes, D. H., Chapin, E. L., Gaztaña, E., Dunlop, J. S., & Ivison, R. J. 2003, MNRAS, 342, 759
Aretxaga, I., Hughes, D. H., & Dunlop, J. S. 2005, MNRAS, 358, 1240
Argon, A. L., et al. 2007, ApJ, 659, 1040
Barger, A. J., Cowie, L. L., Sanders, D. B., Fulton, E., Tanigushi, Y., Sato, Y., Kawara, K., & Okuda, H. 1998, Nature, 394, 248
Barvainis, R., & Antonucci, R. 2005, ApJ, 628, L89
Bertoldi, F., et al. 2000, A&A, 360, 92
Borys, C., et al. 2004, MNRAS, 355, 485
Carilli, C. L., & Yun, M. S. 1999, ApJ, 513, L13
Carilli, C. L., & Yun, M. S. 2000, ApJ, 530, 618
Castangia, P., et al. 2008, A&A, 479, 111
Chapman, S. C., et al. 2003, Nature, 422, 695
Chapman, S. C., et al. 2005, ApJ, 622, 772
Clements, D., et al. 2004, MNRAS, 351, 611
Dannerbauer, H., et al. 2004, ApJ, 606, 664
Darling, J., & Giovanelli, R. 2002, AJ, 124, 100
Gao, Y., et al. 2007, ApJ, 660, L93
Garrett, M. A., Knudsen, K. K., & van der Werf, P. P. 2005, A&A, 431, L21
Genzel, R., & Downes, D. 1979, A&A, 72, 234
Greve, T. R., et al. 2005, MNRAS, 359, 1165
Henkel, C., et al. 2005, A&A, 436, 75
Hughes, D. H., et al. 1998, Nature, 394, 241
Ivison, R. J. 2006, MNRAS, 370, 495
Ivison, R. J., et al. 1998, MNRAS, 298, 583
Ivison, R. J., et al. 2000, ApJ, 542, 27
Ivison, R. J., et al. 2002, MNRAS, 337, 1
Ivison, R. J., et al. 2005, MNRAS, 364, 1025
Ivison, R. J., et al. 2007, MNRAS, 380, 199
Jaffe, D. T., Guesten, R., & Downes, D. 1981, ApJ, 250, 621
Kneib, J.-P., et al. 2004, MNRAS, 349, 1211
Kneib, J.-P., et al. 2005, A&A, 434, 819
Neri, R., et al. 2003, ApJ, 597, L113
Neufeld, D. A., et al. 1994, ApJ, 436, L127
Pérez-Grovas, A. S., Schloerb, F. P., Hughes, D., & Yun, M. 2006, Proc. SPIE, 6267, 1
Smail, I., Ivison, R. J., & Blain, A. W. 1997, ApJ, 490, L5
Smail, I., et al. 2000, ApJ, 528, 612
Solomon, P. M., & Vanden Bout, P. A. 2005, ARA&A, 43, 677
Spiegel, D. N., et al. 2007, ApJS, 170, 377
Townsend, R. H. D., et al. 2001, MNRAS, 328, L17
Wagg, J., et al. 2007, MNRAS, 375, 745
Wang, W.-H., Cowie, L. L., van Saders, J., Barger, A. J., & Williams, J. P. 2007, ApJ, 670, L89
Webb, T. M., et al. 2003, ApJ, 587, 41
Weiß, A., Downes, D., Walter, F., & Henkel, C. 2005, A&A, 440, L45
Wilner, D. J., et al. 1999, ApJ, 117, 1139
Younger, J. D., et al. 2007, ApJ, 671, 1531