Strong enhancement of the critical current at the antiferromagnetic transition in ErNi2B2C single crystals

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Abstract
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We report on transport and magnetization measurements of the critical current density $J_c$ in ErNi$_2$B$_2$C single crystals that show strongly enhanced vortex pinning at the Néel temperature $T_N$ and low applied fields. The height of the observed $J_c$ peak decreases with increasing magnetic field in clear contrast with that of the peak effect found at the upper critical field. Angular transport measurements of $J_c$ revealed the correlated nature of this pinning enhancement, which we attribute to the formation of antiphase boundaries at $T_N$.

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Single quanta of magnetic flux enter a type II superconductor in the form of vortices when it is exposed to a magnetic field $H$ larger than the lower critical field. This is a big setback for applications, because when an electrical current $J$ is applied, vortices move and energy is dissipated. This movement can be arrested by nonsuperconducting defects that pin vortices by lowering the system energy. The interaction between vortices and pinning centers has been studied extensively for decades, focused mainly on the so-called vortex core pinning (caused by the local suppression of the superconducting order parameter). Less explored is the interplay between vortices and magnetic media, which could offer pinning forces superior to those from core pinning. A common difficulty of studying magnetic pinning is the inability of separating the magnetic and core contributions. It is, therefore, of great advantage to study systems where the magnetic transition takes place inside the superconducting phase, allowing one to compare the behavior with and without the magnetically ordered phase.

There are a considerable number of materials in which to study the coexistence between superconductivity and ordered magnetic phases, such as the Chevrel phases, CeCoIn$_5$, and recently iron pnictides. However, the rare earth–nickel–borocarbide family (RENi$_2$B$_2$C, where RE is a rare earth element) has several advantages, namely a relatively high superconducting transition temperature $T_c$ and the tunability of the ratio between its magnetic and superconducting ordering temperatures, which can be changed by using different rare earth elements. The readily available high-purity single crystals allow one to study the effects of magnetic pinning without any significant defect (nonmagnetic) contribution.

Among the borocarbide family the compound with RE = Er has a $T_c$ of 10.5 K and a Néel temperature $T_N$ of 6.0 K. The occurrence of antiferromagnetism directly influences the superconducting properties, as seen in the upper critical field vs temperature curve, where $H_c2(T)$ is slightly suppressed just below $T = T_N$ for $H \parallel c$ and has an inflection point for $H \parallel ab$. This is a consequence of the local moments ordering in the antiferromagnetic phase, as was shown experimentally and theoretically.

In addition to the occurrence of antiferromagnetism, at $T = T_N$ a tetragonal to orthorhombic crystal structure transition takes place. The resulting twin boundaries were shown by Bitter decoration, magneto-optical, and scanning Hall-probe experiments to act as pinning centers at lower temperatures ($T < T_N$). The pinning mechanism was proposed to be caused by a ferromagnetic spin component parallel to the crystallographic $c$ axis localized at the twin boundaries.

Although borocaricides (and particularly the compound with RE = Er) have been extensively studied, few critical current measurements by electrical transport have been performed, with the exception of the seminal work by Gammel et al., exploring the weak ferromagnetism below $T = 2.3$ K. This was shown to lead to an enhanced critical current density $J_c$, which was ascribed to an increase in pinning force due to pair breaking by the ferromagnetism. Canfield et al. pointed out that the data in Ref. extrapolated to $J_c = 0$ at $T_N$, indicating a clear linkage between pinning and the antiferromagnetic order. However, this extrapolation could not be confirmed because transport experiments were only done for $T < 4.2$ K to take advantage of the cooling power of liquid helium.

The lack of transport experiments is to a great extent due to heating problems caused by the applied current being very high because of the large cross section of single crystals. To the best of our knowledge, the present work is the first report both of transport measurements on ErNi$_2$B$_2$C single crystals at temperatures around $T_N$ and of angular $J_c$ measurements performed on this compound.

In this Rapid Communication, we present the results of transport and magnetization measurements on ErNi$_2$B$_2$C single crystals, which revealed a local maximum in $J_c(T)$ at $T = T_N$. We study this large increase in $J_c$ as a function of field strength and orientation, and we find that it occurs only for fields oriented along the crystallographic $c$ axis. Unlike the peak effect near $H_{c2}$, the height of the newly discovered maximum decreases with increasing magnetic field. We rule out pinning from twin boundaries as we are able to determine its angular fingerprint at lower temperatures, which differs from the increase in $J_c$ observed at $T_N$. We attribute the findings to vortex pinning due to antiphase boundaries between antiferromagnetic domains.

Large, homogeneous single crystals of ErNi$_2$B$_2$C were grown using the Ni$_3$B flux growth technique. The samples used for transport measurements were polished mechanically down to a thickness of 40–60 μm parallel to the $c$ axis before being cut to pieces ∼250 μm wide and ∼1.5 mm long.
Sputtered Au contacts allowed us to apply currents along the \( ab \) planes while measuring the voltage across the sample. We observed a slight reduction of \( T_c \) (\(-0.5\) K) due to sample preparation, which is consistent with the previously reported increase in \( T_c \) after annealing.\textsuperscript{31} The small decrease in \( T_c \) does not affect the overall results of the present study; it left both \( T_N \) and the shape of \( H_{c2}(T) \) unchanged.

Transport measurements were performed in three systems characterized by different cooling methods: (i) a Quantum Design, Inc., PPMS with a semiclosed He system, (ii) a variable temperature insert with flowing He, and (iii) with the sample immersed in liquid He. In all systems we used high-precision rotators and the maximum Lorentz force configuration \( (\mathbf{J} \perp \mathbf{H}) \). The critical current density was obtained from \( I-V \) curves using a 1 \( \mu V \) criterion. Measurements of \( R(T) \) and \( R(H) \) gave the upper critical field vs temperature, using a value of \( R \) corresponding to 90\% of the normal conducting resistance. The comparison of the \( J_c \) results from the three systems allowed us to assess the magnitude of heating effects. \( I-V \) curves obtained with the sample immersed in liquid He (which has the best cooling power) served as a reference for the other two systems, allowing us to find the maximum current the PPMS or He flow system could support without heating. All measurements shown are below this threshold.

Magnetization measurements were carried out in a Quantum Design, Inc., MPMS superconducting quantum interference device (SQUID). \( J_c \) was obtained using the Bean critical state model with \( J_c = \frac{\mu_0 |I_m|}{abc \ln(2) \sqrt{b^2 - a^2}} \), where \( a > b, m_1 \) is the irreversible magnetic moment, and \( a, b, \) and \( c \) are the sample length, width, and thickness, respectively.\textsuperscript{32}

In Fig. 1(a) we present a \( J_c \) transport measurement vs temperature at \( \mu_0 H = 0.5 \) T for \( \mathbf{H} \parallel c \). Starting at low temperature, \( J_c(T) \) decreases rapidly as \( T \) increases, and a rough extrapolation of \( J_c(T) \) leads to 0 at \( T = T_N \), the same behavior found in Ref. 24. Our new measurements presented here, however, show that \( J_c \) instead changes abruptly at \( T_N \). The slope of \( J_c(T) \) is much lower at \( T > T_N \) than below, and a peak in \( J_c(T) \) is observed at \( T_N \), corresponding to a twofold increase compared to the value just below and above it. These are the key results of this communication. Finally, at a temperature just above 7 K we observe the well-documented peak effect at \( H_{c2} \).

The maximum at \( T = T_N \) occurs over a range of fields, as seen in the contour plot in Fig. 1(b), obtained from multiple \( J_c(T) \) measurements at different \( \mathbf{H} \parallel c \), where different colors represent different values of \( J_c \). Note that the height of the \( J_c \) peak at \( T_N \) shows the opposite field dependence than that of the one observed near \( H_{c2} \), indicating a different origin of the two maxima. Whereas the peak near \( H_{c2} \) becomes more pronounced as \( H \) is increased, the peak at \( T_N \) decreases with increasing \( H \) until it has essentially vanished at \( \mu_0 H = 0.7 \) T. This leads us to speculate that the peak at \( T_N \) has a magnetic origin, as magnetic pinning tends to be attenuated as the magnetic field modulation of the vortices becomes smaller. Our observations are consistent with the dip found in the dynamic magnetic susceptibility of \( \text{ErNi}_2\text{B}_2\text{C} \) single crystals at \( T = T_N \), measured by Prozorov et al.\textsuperscript{33} using a tunnel-diode resonator technique. They explained their results by a pinning enhancement due to the occurrence of antiferromagnetic order, accompanied by large magnetic fluctuations. The decrease in the peak height with increasing \( H \) also rules out the drastic change in \( H_{c2}(T) \) at \( T_N \) as the origin for the maximum in \( J_c(T) \), since the peak disappears far below \( H_{c2}(T_N) \). The absence of a \( J_c(T) \) peak at \( \theta = 30^\circ \), shown in Fig. 1(a), is further proof that this peak is not associated with the sudden change in \( H_{c2}(T) \) at \( T_N \), since for \( \theta = 30^\circ \) this feature in \( H_{c2}(T) \) is still visible (not shown).

It is clear from the \( J_c(T) \) measurements at \( \theta = 0^\circ \) and \( 30^\circ \) that this effect is strongly angle dependent. Thus, further insight can be gained from the dependence of \( J_c \) on field orientation, since it allows one to distinguish between the superconducting and the magnetic anisotropy.

The obtained angular \( J_c \) curves are reproduced in Fig. 2(a). They reveal rich and complex features. Despite the small change of the absolute value of \( H_{c2} \) with field orientation, we found that \( J_c \) has a pronounced maximum for \( \mathbf{H} \parallel ab \) at all temperatures investigated. This indicates strong pinning for \( \mathbf{H} \parallel ab \), in contrast to what would be expected from the
Figure 2. (Color online) (a) Angular transport $J_c$ measurements at $\mu_0 H = 0.4$ T reveal the narrow temperature range at $T_N$ over which the peak at $H \parallel c$ occurs (obtained in the He flow system). (b) The contour plot of $J_c$ as a function of magnetic field orientation $\theta$ and $T$ shows the location of the peak at $T_N$ and $\theta = 0^\circ$ in the $T$-$\theta$ space (measured in the PPMS).

The dependence of the $T_N$ peak on temperature and field orientation becomes even clearer in the contour plot in Fig. 2(b). It can be seen that $J_c$ is enhanced only over a narrow temperature and angular range near $T = 6$ K and $\theta = 0^\circ$.

The correlated nature of the pinning and its location around $T_N$ lead us to conclude that it is due to dynamic magnetic domains that are created at $T_N$ and annihilated as the temperature decreases further. Domain walls are known to produce planar pinning potentials in ferromagnetic/superconducting hybrids. This poses the interesting question of why the pinning enhancement occurs for $H \parallel c$, while the easy axis of the magnetic moment for ErNi$_2$B$_2$C is the $b$ axis. This could be due to the magnetic moment being flipped to the $c$ axis; it has to be noted, however, that this is highly unlikely because of the high crystalline electric field anisotropy of the local moment sublattice. Another explanation could be vortices bending into the $ab$ planes at the magnetic domain boundaries in order to gain the Zeeman energy. As the vortices move away from the $c$ axis the energy gain gets gradually smaller, until at the $ab$ planes there is no gain since vortices have zero net gain. To shed more light on this topic we are currently investigating the specifics of antiphase boundaries by magnetic force microscopy on ErNi$_2$B$_2$C single crystals. Furthermore, the results presented here motivate transport measurements in other compounds such as TmNi$_2$B$_2$C, which has its easy axis parallel to the $c$ axis. It is worth noting that associated with the tetragonal to orthorhombic transition in ErNi$_2$B$_2$C, strain occurs that could also be responsible for enhanced pinning. Whether this is indeed the case could be answered by investigating a borocarbide which does not undergo a structural transition at $T = T_N$, namely again TmNi$_2$B$_2$C.

In order to further explore the effects at lower fields, which are not accessible to transport measurements due to heating, different samples were taken from the same batch and investigated in the SQUID. This also serves as verification of our transport measurement results. After creating a Bean profile by applying a field $H \parallel c$ larger than $2H^*$ (the field above which screening currents flow in the entire sample volume) at $T = 4.5$ K, the magnetic moment $m$ was measured vs temperature at $H = 0.02$ T. In order to capture the increase in $J_c$ around $T_N$ the Bean profile was generated at different temperatures near $T_N$ before setting $H = 0.02$ T and measuring $m(T)$.

We then compiled an $m(T)$ curve from the different branches, selecting the branch with the highest (absolute)
density. The irreversible magnetic moment then gives the critical current density in ErNi$_2$B$_2$C single crystals at the Néel temperature. By performing the first angular transport measurements we determined that this increase occurs for magnetic fields applied parallel to the $c$ axis, consistent with vortex pinning due to antiphase boundaries between antiferromagnetic domains. This study opens a new avenue to investigate the interaction between superconductivity and different magnetic phases, such as those known in Er, Ho, and Tm borocarbides. In addition, new developments can be expected for other superconductors with coexistence of magnetic phases, such as underdoped copper- and iron-based superconductors.

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FIG. 3. The critical current density vs temperature at $\mu_0H = 0.02$ T (with $\mathbf{H} \parallel c$), obtained from magnetization measurements, exhibits a local maximum at $T = T_N$.

The peak in $J_c$ at $T_N$ and its field dependence are seen in magnetic measurements for both a sample cut and polished to similar dimensions as the one used in transport experiments and a pristine single crystal with no treatment at all. This reinforces that the phenomenon is not sample dependent and can be experimentally verified by different techniques.

In summary, we have observed in transport and magnetization measurements an enhancement of the critical current density in ErNi$_2$B$_2$C single crystals at the Néel temperature. By performing the first angular transport measurements we determined that this increase occurs for magnetic fields applied parallel to the $c$ axis, consistent with vortex pinning due to antiphase boundaries between antiferromagnetic domains. This study opens a new avenue to investigate the interaction between superconductivity and different magnetic phases, such as those known in Er, Ho, and Tm borocarbides. In addition, new developments can be expected for other superconductors with coexistence of magnetic phases, such as underdoped copper- and iron-based superconductors.

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