On the Justification of Viscoelastic Generalized Membrane Equations

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Abstract

We consider a family of linearly viscoelastic shells with thickness $2\varepsilon$, clamped along a portion of their lateral face, all having the same middle surface $S = \theta(\tilde{\omega}) \subset \mathbb{R}^3$, where $\omega \subset \mathbb{R}^2$ is a bounded and connected open set with a Lipschitz-continuous boundary $\gamma$. We show that, if the applied body force density is $O(1)$ with respect to $\varepsilon$ and surface traction density is $O(\varepsilon)$, the solution of the scaled variational problem in curvilinear coordinates, defined over the fixed domain $\Omega = \omega \times (-1,1)$, converges in ad hoc functional spaces as $\varepsilon \to 0$ to a limit $u$. Furthermore, the average $\bar{u}(\varepsilon) = \frac{1}{2} \int_{-1}^{1} u(\varepsilon) dx_3$, converges in an ad hoc space to the unique solution of what we have identified as (scaled) two-dimensional equations of a viscoelastic generalized membrane shell, which includes a long-term memory that takes into account previous deformations. We finally provide convergence results which justify those equations.

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1. Introduction

In solid mechanics, the obtention of models for rods, beams, plates and shells is based on \textit{a priori} hypotheses on the displacement and/or stress fields which, upon substitution in the three-dimensional equilibrium and constitutive equations, lead to useful simplifications. Nevertheless, from both constitutive and geometrical point of views, there is a need to justify the validity of most of the models obtained in this way.

For this reason a considerable effort has been made in the past decades by many authors in order to derive new models and justify the existing ones by using the asymptotic expansion method, whose foundations can be found in \cite{23}. Indeed, the first applied results were obtained with the justification of the linearized theory of plate bending in \cite{11, 14}.

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The existence and uniqueness of solution of elliptic membrane shell equations, can be found in \[13\] and in \[12\]. These two-dimensional models are completely justified with convergence theorems. A complete theory regarding elastic shells can be found in \[10\], where models for elliptic membranes, generalized membranes and flexural shells are presented. It contains a full description of the asymptotic procedure that leads to the corresponding sets of two-dimensional equations. Also, the dynamic case has been studied in \[22, 23, 24\], concerning the justification of dynamic equations for membrane, flexural and Koiter shells.

A large number of real problems had made it necessary the study of new models which could take into account effects such as hardening and memory of the material. An example of these are the viscoelasticity models (see \[13, 21, 28\]). Many authors have contributed to the nowadays knowledge of this sort of problems, providing justified models and results. Indeed, we can find examples in the literature as \[17, 18, 19, 20, 26, 29\] and in the references therein, a variety of models for problems concerning the viscoelastic behaviour of the material. In particular, there exist studies of the behaviour of viscoelastic plates as in \[3, 5\], where models for von Kármán plates are analysed. In some of these works, we can find analysis of the influence of short or long term memory in the equations modelling a problem. These terms take into account previous deformations of the body, hence, they are commonly presented in some viscoelastic problems. For instance, on one hand, we can find in \[1\] models including a short term memory presented by a system of integro-differential and pseudoparabolic equations describing large deflections on a viscoelastic plate. On the other hand, in \[27\] a long term memory is considered on the study of the asymptotic behaviour of the solution of a von Kármán plate when the time variable tends to infinity. Also, in the reference \[2\], the authors study the effects of great deflections in thin plates covering both short and long term memory cases. Concerning viscoelastic shell problems, in \[32\] we can find different kind of studies where the authors also remark the viscoelastic property of the material of a shell. For the problems dealing with the shell-type equations, there exists a very limited amount of results available, for instance, \[4\] where the authors present a model for a dynamic contact problem where a short memory (Kelvin-Voigt) material is considered. Particularly remarkable is the increasing number of studies of viscoelastic shells problems in order to reproduce the complex behaviour of tissues in the field of biomechanics. For example, in \[16\] the difficulties of this kind of problems are detailed and even though an one-dimensional model is derived for modelling a vessel wall, the author comments the possibility of considering two-dimensional models with a shell-type description and a viscoelastic constitutive law. In this direction, to our knowledge, in \[8\] we gave the first steps towards the justification of existing models of viscoelastic shells and the finding of new ones. By using the asymptotic expansion method, we found a rich variety of cases, depending on the geometry of the middle surface, the boundary conditions and the order of the applied forces. The most remarkable feature was that from the asymptotic analysis of the three-dimensional problems which included a short term memory represented by a time derivative, a long term memory arised in the two-dimensional limit problems, represented by an integral with respect to the time variable. This fact, agreed with previous asymptotic analysis of viscoelastic rods in \[30, 31\] where an analogous behaviour was presented as well.

In \[9, 7\] we justified the equations of a viscoelastic membrane shell where the surface
$S$ is elliptic and the boundary condition of place is considered in the whole lateral face of the shell. Therefore, in this paper the main aim is to justify the remaining cases in the group of viscoelastic membrane cases, known as the viscoelastic generalized membrane shell equations. In such a group, we shall distinguish two kinds of membranes, that will be detailed in following sections. As an example, we present one of the set of equations that will be justified in this paper:

**Problem 1.1.** Find $\xi^e(t, \cdot) \in V^\#_M(\omega)$ $\forall t \in [0, T]$ such that,

$$
B^{\#}_{M,t}(\xi^e, \eta) = L^{\#}_{M,t}(\eta) \forall \eta \in V^\#_M(\omega), \text{ a.e. } t \in (0, T),
$$

$$
\xi^e(0, \cdot) = \xi^e_0(\cdot),
$$

where $B^{\#}_{M,t}$ and $L^{\#}_{M,t}$ are the unique continuous extensions from $W^{1,2}(0, T, V(\omega))$ to $W^{1,2}(0, T, V^\#_M(\omega))$ and from $V(\omega)$ to $V^\#_M(\omega)$ of the functions $B^{\#}_{M,t} : W^{1,2}(0, T, V(\omega)) \times V(\omega) \rightarrow \mathbb{R}$ and $L^{\#}_{M,t} : V(\omega) \rightarrow \mathbb{R}$, respectively, defined by

$$
B^{\#}_{M,t}(\xi^e, \eta) := \varepsilon \int_\omega \alpha^{\beta\sigma\tau, e}\gamma_{\sigma\tau}(\xi^e(t))\gamma_{\alpha\beta}(\eta)\sqrt{a}dy + \varepsilon \int_\omega b^{\alpha\beta\sigma\tau, e}\gamma_{\sigma\tau}(\xi^e(t))\gamma_{\alpha\beta}(\eta)\sqrt{a}dy
$$

$$
- \varepsilon \int_0^t e^{-k(t-s)} \int_\omega c^{\alpha\beta\sigma\tau, e}\gamma_{\sigma\tau}(\xi^e(s))\gamma_{\alpha\beta}(\eta)\sqrt{a}dyds,
$$

$$
L^{\#}_{M,t}(\eta) := \int_\omega \varphi^{\alpha\beta, e}(t)\gamma_{\alpha\beta}(\eta)\sqrt{a}dy,
$$

where $V^\#_M(\omega)$ is an ad hoc abstract space in which this problem is well posed (defined in Section 2), $\alpha^{\beta\sigma\tau, e}$, $b^{\alpha\beta\sigma\tau, e}$ and $c^{\alpha\beta\sigma\tau, e}$ denote the re-scaled versions of the contravariant components of the two-dimensional fourth order tensors that we shall recall later (6.1)–(6.3), $\varphi^{\alpha\beta, e}$ is a de-scaled version of the real function defined in (6.9) derived from the convergence study which depends on the applied forces and $\gamma_{\alpha\beta}(\cdot)$ denote the covariant components of the linearized change of metric tensor associated with a displacement field of the surface $S$, which we will recall in (4.4).

We prove that the scaled three-dimensional unknown, $u(\varepsilon)$, converges as the small parameter $\varepsilon$ tends to zero in an ad hoc functional space and its transversal average converges to $\xi^e$, the unique solution of the two-dimensional associated problem. Moreover, unlike the viscoelastic elliptic membrane shells, the limit of the scaled three-dimensional unknown $u(\varepsilon)$ is not necessary independent of $x_3$, however we find that that $\partial_3 u(\varepsilon) \rightarrow 0$ in $W^{1,2}(0, T, L^2(\Omega))$.

We will follow the notation and style of [10], where the linear elastic shells are studied. For this reason, we shall reference auxiliary results which apply in the same manner to the viscoelastic case. One of the major differences with respect to previous works in elasticity, consists on time dependence, that will lead to ordinary differential equations that need to be solved in order to find the zeroth-order approach of the solution.

The structure of the paper is the following: in Section 2 we shall recall the viscoelastic problem in Cartesian coordinates and then, considering the problem for a family of viscoelastic shells of thickness $2\varepsilon$, we formulate the problem in curvilinear coordinates. In Section 3
we will use a projection map into a reference domain, we will introduce the scaled unknowns and forces and the assumptions on coefficients. In Section 4 we recall some technical results which will be needed in what follows. In Section 5 we shall study the completion spaces that will lead to well posed problems for the viscoelastic membrane shell equations. Then, we will introduce an assumption on the applied forces, needed in the convergence analysis. In Section 6 we enunciate the two-dimensional equations for a viscoelastic generalized membrane shell and we present the convergence results when the small parameter $\varepsilon$ tends to zero, which is the main result of this paper. Then, we present the convergence results in terms of de-scaled unknowns. In Section 7 we shall present some conclusions, including a comparison between the viscoelastic models and the elastic case studied in [10] and announce the convergence results regarding other cases in forthcoming papers.

2. The three-dimensional linearly viscoelastic shell problem

We denote by $\mathbb{S}^d$, where $d = 2, 3$ in practice, the space of second-order symmetric tensors on $\mathbb{R}^d$, while “ $\cdot , \cdot $” will represent the inner product and $| \cdot |$ the usual norm in $\mathbb{S}^d$ and $\mathbb{R}^d$. In what follows, unless the contrary is explicitly written, we will use summation convention on repeated indices. Moreover, Latin indices $i, j, k, l, \ldots$, take their values in the set $\{1, 2, 3\}$, whereas Greek indices $\alpha, \beta, \sigma, \tau, \ldots$, do it in the set $\{1, 2\}$. Also, we use standard notation for the Lebesgue and Sobolev spaces. Also, for a time dependent function $u$, we denote $\dot{u}$ the first derivative of $u$ with respect to the time variable. Recall that “ $\rightarrow$ ” denotes strong convergence, while “ $\rightharpoonup$ ” denotes weak convergence.

Let $\Omega^*$ be a domain of $\mathbb{R}^3$, with a Lipschitz-continuous boundary $\Gamma^* = \partial \Omega^*$. Let $x^*(x_i^*)$ be a generic point of its closure $\bar{\Omega}^*$ and let $\partial_i^*$ denote the partial derivative with respect to $x_i^*$. Let $dx^*$ denote the volume element in $\Omega^*$, $d\Gamma^*$ denote the area element along $\Gamma^*$ and $n^*$ denote the unit outer normal vector along $\Gamma^*$. Finally, let $\Gamma_0^*$ and $\Gamma_1^*$ be subsets of $\Gamma^*$ such that $\text{meas}(\Gamma_0^*) > 0$ and $\Gamma_0^* \cap \Gamma_1^* = \emptyset$.

The set $\Omega^*$ is the region occupied by a deformable body in the absence of applied forces. We assume that this body is made of a Kelvin-Voigt viscoelastic material, which is homogeneous and isotropic, so that the material is characterized by its Lamé coefficients $\lambda \geq 0, \mu > 0$ and its viscosity coefficients, $\theta \geq 0, \rho \geq 0$ (see for instance [15, 21, 33]).

Let $T > 0$ be the time period of observation. Under the effect of applied forces, the body is deformed and we denote by $u_i^* : [0, T] \times \bar{\Omega}^* \to \mathbb{R}^3$ the Cartesian components of the displacements field, defined as $u^* := u_i^* e_i^* : [0, T] \times \bar{\Omega}^* \to \mathbb{R}^3$, where $\{e_i^*\}$ denotes the Euclidean canonical basis in $\mathbb{R}^3$. Moreover, we consider that the displacement field vanishes on the set $\Gamma_0^*$. Hence, the displacements field $u^* = (u_i^*) : [0, T] \times \Omega^* \to \mathbb{R}^3$ is solution of the following three-dimensional problem in Cartesian coordinates.
Problem 2.1. Find $u^* = (u^*_i) : [0, T] \times \Omega^* \rightarrow \mathbb{R}^3$ such that,

\begin{align*}
-\partial^*_j \sigma^{ij,*}(u^*) &= f^{i,*} \quad \text{in } \Omega^*, \quad (2.1) \\
u^*_i &= 0 \quad \text{on } \Gamma^*_0, \quad (2.2) \\
\sigma^{ij,*}(u^*) n_j^* &= h^{i,*} \quad \text{on } \Gamma^*_1, \quad (2.3) \\
u^*(0, \cdot) &= u^*_0 \quad \text{in } \Omega^*, \quad (2.4)
\end{align*}

where the functions

\[
\sigma^{ij,*}(u^*) := A^{ijkl,*} e^{kl,*}_{ij}(u^*) + B^{ijkl,*} \dot{e}^{kl,*}_{ij}(\dot{u}^*),
\]

are the components of the linearized stress tensor field and where the functions

\[
A^{ijkl,*} := \lambda \delta^{ij} \delta^{kl} + \mu \left( \delta^{ik} \delta^{jl} + \delta^{il} \delta^{jk} \right),
\]

\[
B^{ijkl,*} := \theta \delta^{ij} \delta^{kl} + \rho \left( \delta^{ik} \delta^{jl} + \delta^{il} \delta^{jk} \right),
\]

are the components of the three-dimensional elasticity and viscosity fourth order tensors, respectively, and

\[
e^{ij,*}(u^*) := \frac{1}{2} (\partial^*_j u^*_i + \partial^*_i u^*_j),
\]

designates the components of the linearized strain tensor associated with the displacement field $u^*$ of the set $\Omega^*$.

We now proceed to describe the equations in Problem 2.1. Expression (2.1) is the equilibrium equation, where $f^{i,*}$ are the components of the volumic force densities. The equality (2.2) is the Dirichlet condition of place, (2.3) is the Neumann condition, where $h^{i,*}$ are the components of surface force densities and (2.4) is the initial condition, where $u^*_0$ denotes the initial displacements.

Note that, for the sake of briefness, we omit the explicit dependence on the space and time variables when there is no ambiguity. Let us define the space of admissible unknowns,

\[
V(\Omega^*) = \{ v^* = (v^*_i) \in [H^1(\Omega^*)]^3; v^*_i = 0 \text{ on } \Gamma^*_0 \}. 
\]

Therefore, assuming enough regularity, the unknown $u^* = (u^*_i)$ satisfies the following variational problem in Cartesian coordinates:

Problem 2.2. Find $u^* = (u^*_i) : [0, T] \times \Omega^* \rightarrow \mathbb{R}^3$ such that,

\[
\int_{\Omega^*} A^{ijkl,*} e^{kl,*}_{ij}(u^*(t)) e^{ij,*}(v^*) dx^* + \int_{\Omega^*} B^{ijkl,*} \dot{e}^{kl,*}_{ij}(\dot{u}^*(t)) e^{ij,*}(v^*) dx^* \\
\quad = \int_{\Omega^*} f^{i,*}(t) v^*_i dx^* + \int_{\Gamma^*_1} h^{i,*}(t) v^*_i d\Gamma^* \quad \forall v^* \in V(\Omega^*), \ a.e. \ in \ (0, T), \\
u^*(0, \cdot) = u^*_0(\cdot). 
\]
Let us consider that $\Omega^*$ is a viscoelastic shell of thickness $2\varepsilon$. Now, we shall express the equations of the Problem 2.2 in terms of curvilinear coordinates. Let $\omega$ be a domain of $\mathbb{R}^2$, with a Lipschitz-continuous boundary $\gamma = \partial \omega$. Let $y = (y_a)$ be a generic point of its closure $\bar{\omega}$ and let $\partial_a$ denote the partial derivative with respect to $y_a$.

Let $\theta \in C^2(\bar{\omega}; \mathbb{R}^3)$ be an injective mapping such that the two vectors $a_\alpha(y) := \partial_\alpha \theta(y)$ are linearly independent. These vectors form the covariant basis of the tangent plane to the surface $S := \theta(\bar{\omega})$ at the point $\theta(y) = y^*$. We can consider the two vectors $a^\alpha(y)$ of the same tangent plane defined by the relations $a^\alpha(y) \cdot a_\beta(y) = \delta^\alpha_\beta$, that constitute the contravariant basis. We define the unit vector,

$$
a_3(y) = a^3(y) := \frac{a_1(y) \wedge a_2(y)}{|a_1(y) \wedge a_2(y)|},
$$

(2.5)

normal vector to $S$ at the point $\theta(y) = y^*$, where $\wedge$ denotes vector product in $\mathbb{R}^3$.

We can define the first fundamental form, given as metric tensor, in covariant or contravariant components, respectively, by

$$a_{\alpha\beta} := a_\alpha \cdot a_\beta, \quad a^{\alpha\beta} := a^\alpha \cdot a^\beta,$$

the second fundamental form, given as curvature tensor, in covariant or mixed components, respectively, by

$$b_{\alpha\beta} := a_3 \cdot \partial_\beta a_\alpha, \quad b_\alpha^\beta := a_\beta^\sigma b_{\sigma\alpha},$$

and the Christoffel symbols of the surface $S$ by

$$\Gamma^\sigma_{\alpha\beta} := a^\sigma \cdot \partial_\beta a_\alpha.$$

The area element along $S$ is $\sqrt{a} dy = dy^*$ where

$$a := \det(a_{\alpha\beta}).$$

(2.6)

Let $\gamma_0$ be a subset of $\gamma$, such that $\text{meas}(\gamma_0) > 0$. For each $\varepsilon > 0$, we define the three-dimensional domain $\Omega^\varepsilon := \omega \times (-\varepsilon, \varepsilon)$ and its boundary $\Gamma^\varepsilon = \partial \Omega^\varepsilon$. We also define the following parts of the boundary,

$$\Gamma^\varepsilon_+ := \omega \times \{\varepsilon\}, \quad \Gamma^\varepsilon_- := \omega \times \{-\varepsilon\}, \quad \Gamma^\varepsilon_0 := \gamma_0 \times [-\varepsilon, \varepsilon].$$

Let $x^\varepsilon = (x^\varepsilon_1)$ be a generic point of $\bar{\Omega}^\varepsilon$ and let $\partial^\varepsilon_i$ denote the partial derivative with respect to $x^\varepsilon_i$. Note that $x^\varepsilon_\alpha = y_\alpha$ and $\partial^\varepsilon_\alpha = \partial_\alpha$. Let $\Theta : \bar{\Omega}^\varepsilon \to \mathbb{R}^3$ be the mapping defined by

$$\Theta(x^\varepsilon) := \theta(y) + x^\varepsilon_3 a_3(y) \forall x^\varepsilon = (y, x^\varepsilon_3) \in \bar{\Omega}^\varepsilon.$$

(2.7)

The next theorem shows that if the injective mapping $\theta : \bar{\omega} \to \mathbb{R}^3$ is smooth enough, the mapping $\Theta : \bar{\Omega}^\varepsilon \to \mathbb{R}^3$ is also injective for $\varepsilon > 0$ small enough (see Theorem 3.1-1, [10]).
Theorem 2.3. Let $\omega$ be a domain in $\mathbb{R}^2$. Let $\theta \in C^2(\bar{\omega}; \mathbb{R}^3)$ be an injective mapping such that the two vectors $a_\alpha = \partial_\alpha \theta$ are linearly independent at all points of $\bar{\omega}$ and let $a_3$, defined in $\mathbb{R}^3$. Then there exists $\varepsilon_0 > 0$ such that the mapping $\Theta : \Omega_0 \to \mathbb{R}^3$ defined by

$$
\Theta(y, x_3) := \theta(y) + x_3 a_3(y) \forall (y, x_3) \in \Omega_0, \text{ where } \Omega_0 := \omega \times (-\varepsilon_0, \varepsilon_0),
$$

is a $C^1$--diffeomorphism from $\Omega_0$ onto $\Theta(\Omega_0)$ and $\det(g_1, g_2, g_3) > 0$ in $\Omega_0$, where $g_i := \partial_i \Theta$.

For each $\varepsilon$, $0 < \varepsilon \leq \varepsilon_0$, the set $\Theta(\Omega^\varepsilon) = \bar{\Omega}^\varepsilon$ is the reference configuration of a viscoelastic shell, with middle surface $S = \theta(\bar{\omega})$ and thickness $2\varepsilon > 0$. Furthermore for $\varepsilon > 0$, $g^\varepsilon_i(x^\varepsilon) := \partial_i^{\varepsilon} \Theta(x^\varepsilon)$ are linearly independent and the mapping $\Theta : \Omega^\varepsilon \to \mathbb{R}^3$ is injective for all $\varepsilon$, $0 < \varepsilon \leq \varepsilon_0$, as a consequence of injectivity of the mapping $\theta$. Hence, the three vectors $g^\varepsilon_i(x^\varepsilon)$ form the covariant basis of the tangent space at the point $x^\varepsilon = \Theta(x^\varepsilon)$ and $g^{i,\varepsilon}(x^\varepsilon)$ defined by the relations $g^{i,\varepsilon} \cdot g^{j,\varepsilon} = \delta^i_j$ form the contravariant basis at the point $x^\varepsilon = \Theta(x^\varepsilon)$. We define the metric tensor, in covariant or contravariant components, respectively, by

$$
g^\varepsilon_{ij} := g^\varepsilon_i \cdot g^\varepsilon_j, \quad g^{i,\varepsilon} \cdot g^{j,\varepsilon},
$$

and Christoffel symbols by

$$
\Gamma^{p,\varepsilon}_{ij} := g^{p,\varepsilon} \cdot \partial_i^{\varepsilon} g^\varepsilon_j. \quad (2.8)
$$

The volume element in the set $\Theta(\Omega^\varepsilon) = \bar{\Omega}^\varepsilon$ is $\sqrt{g^\varepsilon} dx^\varepsilon = dx^s$ and the surface element in $\Theta(\Gamma^\varepsilon) = \Gamma^s$ is $\sqrt{g^\varepsilon} d\Gamma^s = d\Gamma^s$ where

$$
g^\varepsilon := \det(g^\varepsilon_{ij}). \quad (2.9)
$$

Therefore, for a field $v^s$ defined in $\Theta(\bar{\Omega}^\varepsilon) = \bar{\Omega}^s$, we define its covariant curvilinear coordinates $v^s_i$ by

$$
v^s(x^s) = v^s_i(x^s) e^i =: v^s_i(x^s) g^i(x^\varepsilon), \text{ with } x^s = \Theta(x^\varepsilon).
$$

Besides, we denote by $u^\varepsilon_i : [0, T] \times \Omega^\varepsilon \to \mathbb{R}^3$ the covariant components of the displacements field, that is $U^\varepsilon := u^\varepsilon_i g^i, \quad [0, T] \times \Omega^\varepsilon \to \mathbb{R}^3$. For simplicity, we define the vector field $u^\varepsilon = (u^\varepsilon_i) : [0, T] \times \Omega^\varepsilon \to \mathbb{R}^3$ which will be denoted vector of unknowns.

Recall that we assumed that the shell is subjected to a boundary condition of place; in particular that the displacements field vanishes in $\Theta(\Gamma^\varepsilon_0) = \Gamma^s_0$.

Accordingly, let us define the space of admissible unknowns,

$$
V(\Omega^\varepsilon) = \{ v^\varepsilon = (v^\varepsilon_i) \in [H^1(\Omega^\varepsilon)]^3; v^\varepsilon = 0 \text{ on } \Gamma^s_0 \}.
$$

This is a real Hilbert space with the induced inner product of $[H^1(\Omega^\varepsilon)]^3$. The corresponding norm is denoted by $\| \cdot \|_{1,\Omega^\varepsilon}$.

Therefore, we can find the expression of the Problem 2.2 in curvilinear coordinates (see [10] for details). Hence, the “displacements” field $u^\varepsilon = (u^\varepsilon_i)$ verifies the following variational problem of a three-dimensional viscoelastic shell in curvilinear coordinates:
Problem 2.4. Find \( \mathbf{u}^\varepsilon = (u_i^\varepsilon) : [0, T] \times \Omega^\varepsilon \to \mathbb{R}^3 \) such that,

\[
\mathbf{u}^\varepsilon(t, \cdot) \in V(\Omega^\varepsilon) \quad \forall \, t \in [0, T],
\]

\[
\int_{\Omega^\varepsilon} A^{ijkl,\varepsilon} e^{\varepsilon}_{k||l}(\mathbf{u}^\varepsilon(t)) e^{\varepsilon}_{i||j}(\mathbf{v}) \sqrt{g^\varepsilon} \, dx^\varepsilon + \int_{\Omega^\varepsilon} B^{ijkl,\varepsilon} e^{\varepsilon}_{k||l}(\mathbf{u}^\varepsilon(t)) e^{\varepsilon}_{i||j}(\mathbf{v}) \sqrt{g^\varepsilon} \, dx^\varepsilon
\]

\[
= \int_{\Omega^\varepsilon} f^{i,\varepsilon}(t) v_i^\varepsilon \sqrt{g^\varepsilon} \, dx^\varepsilon + \int_{\Gamma_{\varepsilon}^+ \cup \Gamma_{\varepsilon}^-} h^{i,\varepsilon}(t) v_i^\varepsilon \sqrt{g^\varepsilon} \, d\Gamma^\varepsilon \quad \forall \mathbf{v}^\varepsilon \in V(\Omega^\varepsilon), \ \text{a.e. in } (0, T),
\]

\[
\mathbf{u}^\varepsilon(0, \cdot) = \mathbf{u}_0^\varepsilon(\cdot),
\]

where the functions

\[
A^{ijkl,\varepsilon} := \lambda g^{ij,\varepsilon} g^{kl,\varepsilon} + \mu (g^{ik,\varepsilon} g^{jl,\varepsilon} + g^{il,\varepsilon} g^{jk,\varepsilon}),
\]

\[
B^{ijkl,\varepsilon} := \theta g^{ij,\varepsilon} g^{kl,\varepsilon} + \frac{\mu}{2} (g^{ik,\varepsilon} g^{jl,\varepsilon} + g^{il,\varepsilon} g^{jk,\varepsilon}),
\]

are the contravariant components of the three-dimensional elasticity and viscosity tensors, respectively. We assume that the Lamé coefficients \( \lambda \geq 0, \mu > 0 \) and the viscosity coefficients \( \theta \geq 0, \rho \geq 0 \) are all independent of \( \varepsilon \). Moreover, the terms

\[
e^{\varepsilon}_{i||j}(\mathbf{u}^\varepsilon) := \frac{1}{2} (u^\varepsilon_{i||j} + u^\varepsilon_{j||i}) = \frac{1}{2} (\partial_j u_i^\varepsilon + \partial_i u_j^\varepsilon) - \Gamma_{ij}^{p,\varepsilon} u^\varepsilon_p,
\]

de designate the covariant components of the linearized strain tensor associated with the displacement field \( \mathbf{U}^\varepsilon \) of the set \( \Theta(\bar{\Omega}^\varepsilon) \). Moreover, \( f^{i,\varepsilon} \) denotes the contravariant components of the volumic force densities, \( h^{i,\varepsilon} \) denotes contravariant components of surface force densities and \( \mathbf{u}_0^\varepsilon \) denotes the initial " displacements " (actually, the initial displacement is \( \mathbf{U}_0^\varepsilon := (u_0^\varepsilon), g^{i,\varepsilon} \)).

Note that the following additional relations are satisfied,

\[
\Gamma_{a3}^{3,\varepsilon} = \Gamma_{33}^{p,\varepsilon} = 0 \text{ in } \bar{\Omega}^\varepsilon,
\]

\[
A^{033,\varepsilon} = A^{033,\varepsilon} = B^{033,\varepsilon} = B^{033,\varepsilon} = 0 \text{ in } \bar{\Omega}^\varepsilon,
\]

as a consequence of the definition of \( \Theta \) in (2.7).

The existence and uniqueness of solution of the Problem 2.4 for \( \varepsilon > 0 \) small enough, established in the following theorem, was proved in [8].

Theorem 2.5. Let \( \Omega^\varepsilon \) be a domain in \( \mathbb{R}^3 \) defined previously in this section and let \( \Theta \) be a \( \mathcal{C}^2 \)-diffeomorphism of \( \bar{\Omega}^\varepsilon \) in its image \( \Theta(\bar{\Omega}^\varepsilon) \), such that the three vectors \( \mathbf{g}^{\varepsilon}_i(\mathbf{x}) = \partial_i \Theta(\mathbf{x}^\varepsilon) \) are linearly independent for all \( \mathbf{x}^\varepsilon \in \bar{\Omega}^\varepsilon \). Let \( \Gamma_0^\varepsilon \) be a \( d\Gamma^\varepsilon \)-measurable subset of \( \gamma \times [-\varepsilon, \varepsilon] \) such that \( \text{meas}(\Gamma_0^\varepsilon) > 0 \). Let \( f^{i,\varepsilon} \in L^2(0, T; L^2(\Omega^\varepsilon)) \), \( h^{i,\varepsilon} \in L^2(0, T; L^2(\Gamma_i^\varepsilon)) \), where \( \Gamma_i^\varepsilon := \Gamma_i^+ \cup \Gamma_i^- \). Let \( \mathbf{u}_0^\varepsilon \in V(\bar{\Omega}^\varepsilon) \). Then, there exists a unique solution \( \mathbf{u}^\varepsilon = (u_i^\varepsilon) : [0, T] \times \bar{\Omega}^\varepsilon \to \mathbb{R}^3 \) satisfying the Problem 2.4. Moreover, \( \mathbf{u}^\varepsilon \in W^{1,2}(0, T; V(\Omega^\varepsilon)) \). In addition to that, if \( \dot{f}^{i,\varepsilon} \in L^2(0, T; L^2(\Omega^\varepsilon)) \), \( \dot{h}^{i,\varepsilon} \in L^2(0, T; L^2(\Gamma_i^\varepsilon)) \), then \( \mathbf{u}^\varepsilon \in W^{2,2}(0, T; V(\bar{\Omega}^\varepsilon)) \).
3. The scaled three-dimensional shell problem

For convenience, we consider a reference domain independent of the small parameter $\varepsilon$. Hence, let us define the three-dimensional domain $\Omega := \omega \times (-1, 1)$ and its boundary $\Gamma = \partial \Omega$. We also define the following parts of the boundary,

$$\Gamma_+ := \omega \times \{1\}, \quad \Gamma_- := \omega \times \{-1\}, \quad \Gamma_0 := \gamma_0 \times [-1, 1].$$

Let $x = (x_1, x_2, x_3)$ be a generic point in $\bar{\Omega}$ and we consider the notation $\partial_i$ for the partial derivative with respect to $x_i$. We define the following projection map,

$$\pi^\varepsilon : x = (x_1, x_2, x_3) \in \bar{\Omega} \rightarrow \pi^\varepsilon(x) = x^\varepsilon = (x_1^\varepsilon, x_2^\varepsilon, x_3^\varepsilon) = (x_1, x_2, \varepsilon x_3) \in \bar{\Omega}^\varepsilon,$$

hence, $\partial_\alpha^\varepsilon = \partial_\alpha$ and $\partial_3^\varepsilon = \frac{1}{\varepsilon} \partial_3$. We consider the scaled unknown $u(\varepsilon) = (u_i(\varepsilon)) : [0, T] \times \bar{\Omega} \rightarrow \mathbb{R}^3$ and the scaled vector fields $v = (v_i) : \bar{\Omega} \rightarrow \mathbb{R}^3$ defined as

$$u_i^\varepsilon(t, x^\varepsilon) := u_i(\varepsilon)(t, x) \text{ and } v_i^\varepsilon(x^\varepsilon) := v_i(x) \forall x^\varepsilon = \pi^\varepsilon(x) \in \bar{\Omega}^\varepsilon, \forall t \in [0, T].$$

Also, let the functions, $\Gamma_{ij}^{\varepsilon}, g^\varepsilon, A^{ijkl,\varepsilon}, B^{ijkl,\varepsilon}$ defined in (2.8), (2.9), (2.11) and (2.12), be associated with the functions $\Gamma_{ij}^{\varepsilon}(\varepsilon), g(\varepsilon), A^{ijkl}(\varepsilon), B^{ijkl}(\varepsilon)$ defined by

$$\Gamma_{ij}^{\varepsilon}(\varepsilon)(x) := \Gamma_{ij}^{\varepsilon}(x^\varepsilon), \quad \text{(3.1)}$$
$$g(\varepsilon)(x) := g^\varepsilon(x^\varepsilon), \quad \text{(3.2)}$$
$$A^{ijkl}(\varepsilon)(x) := A^{ijkl,\varepsilon}(x^\varepsilon), \quad \text{(3.3)}$$
$$B^{ijkl}(\varepsilon)(x) := B^{ijkl,\varepsilon}(x^\varepsilon), \quad \text{(3.4)}$$

for all $x^\varepsilon = \pi^\varepsilon(x) \in \bar{\Omega}^\varepsilon$. For all $v = (v_i) \in [H^1(\Omega)]^3$, let there be associated the scaled linearized strains components $e_{i||j}(\varepsilon; v) \in L^2(\Omega)$, defined by

$$e_{\alpha||\beta}(\varepsilon; v) := \frac{1}{2} (\partial_\beta v_\alpha + \partial_\alpha v_\beta) - \Gamma_{\alpha\beta}(\varepsilon)v_p, \quad \text{(3.5)}$$
$$e_{\alpha||3}(\varepsilon; v) := \frac{1}{2} \left( \frac{1}{\varepsilon} \partial_3 v_\alpha + \partial_\alpha v_3 \right) - \Gamma_{\alpha3}(\varepsilon)v_p, \quad \text{(3.6)}$$
$$e_{3||3}(\varepsilon; v) := \frac{1}{\varepsilon} \partial_3 v_3. \quad \text{(3.7)}$$

Note that with these definitions it is verified that

$$e_{i||j}^\varepsilon(\pi^\varepsilon(x)) = e_{i||j}(\varepsilon; v)(x) \forall x \in \Omega.$$

**Remark 3.1.** The functions $\Gamma_{ij}^{\varepsilon}(\varepsilon), g(\varepsilon), A^{ijkl}(\varepsilon), B^{ijkl}(\varepsilon)$ converge in $C^0(\bar{\Omega})$ when $\varepsilon$ tends to zero.

**Remark 3.2.** When we consider $\varepsilon = 0$ the functions will be defined with respect to $y \in \bar{\omega}$.

We shall distinguish the three-dimensional Christoffel symbols from the two-dimensional ones by using $\Gamma_{\alpha\beta}^{\varepsilon}(\varepsilon)$ and $\Gamma_{\alpha\beta}^\sigma$, respectively.
The next result is an adaptation of (b) in Theorem 3.3-2, [10] to the viscoelastic case. We will study the asymptotic behaviour of the scaled contravariant components \( A_{ijkl}(\varepsilon) \), \( B_{ijkl}(\varepsilon) \) of the three-dimensional elasticity and viscosity tensors defined in (3.3)–(3.4), as \( \varepsilon \to 0 \). We show their uniform positive definiteness not only with respect to \( \mathbf{x} \in \overline{\Omega} \), but also with respect to \( \varepsilon, 0 < \varepsilon \leq \varepsilon_0 \). Finally, their limits are functions of \( \mathbf{y} \in \overline{\omega} \) only, that is, independent of the transversal variable \( x_3 \).

**Theorem 3.3.** Let \( \omega \) be a domain in \( \mathbb{R}^2 \) and let \( \mathbf{\theta} \in C^2(\overline{\omega}; \mathbb{R}^3) \) be an injective mapping such that the two vectors \( \mathbf{a}_\alpha = \partial_\alpha \mathbf{\theta} \) are linearly independent at all points of \( \overline{\omega} \), let \( a^{\alpha\beta} \) denote the contravariant components of the metric tensor of \( S = \mathbf{\theta}(\overline{\omega}) \). In addition to that, let the other assumptions on the mapping \( \mathbf{\theta} \) and the definition of \( \varepsilon_0 \) be as in Theorem 2.3. The contravariant components \( A_{ijkl}(\varepsilon) \), \( B_{ijkl}(\varepsilon) \) of the scaled three-dimensional elasticity and viscosity tensors, respectively, defined in (3.3)–(3.4) satisfy

\[
A_{ijkl}(\varepsilon) = A_{ijkl}(0) + O(\varepsilon) \quad \text{and} \quad A_{\alpha\beta\sigma\tau}(\varepsilon) = A_{\alpha\beta\sigma\tau}(0) = 0,
\]
\[
B_{ijkl}(\varepsilon) = B_{ijkl}(0) + O(\varepsilon) \quad \text{and} \quad B_{\alpha\beta\sigma\tau}(\varepsilon) = B_{\alpha\beta\sigma\tau}(0) = 0,
\]
for all \( \varepsilon, 0 < \varepsilon \leq \varepsilon_0 \), and

\[
A_{\alpha\beta\tau}(0) = \lambda a^{\alpha\beta} a^{\sigma\tau} + \mu (a^{\alpha\sigma} a^{\beta\tau} + a^{\alpha\tau} a^{\beta\sigma}), \quad A_{\alpha\beta\sigma\tau}(0) = \lambda a^{\alpha\beta},
\]
\[
A_{\alpha\beta\sigma}(0) = \mu a^{\alpha\sigma}, \quad A_{3333}(0) = \lambda + 2\mu,
\]
\[
B_{\alpha\beta\tau}(0) = \theta a^{\alpha\beta} a^{\sigma\tau} + \frac{\rho}{2} (a^{\alpha\sigma} a^{\beta\tau} + a^{\alpha\tau} a^{\beta\sigma}), \quad B_{\alpha\beta\sigma\tau}(0) = \theta a^{\alpha\beta},
\]
\[
B_{\alpha\beta\sigma}(0) = \frac{\rho}{2} a^{\alpha\sigma}, \quad B_{3333}(0) = \theta + \rho,
\]
\[
B_{\alpha\beta\sigma\tau}(0) = B_{\alpha\beta\sigma}(0) = 0.
\]

Moreover, there exist two constants \( C_e > 0 \) and \( C_v > 0 \), independent of the variables and \( \varepsilon \), such that

\[
\sum_{i,j} |t_{ij}|^2 \leq C_e A_{ijkl}(\varepsilon)(\mathbf{x}) t_{kl} t_{ij}, \quad (3.8)
\]
\[
\sum_{i,j} |t_{ij}|^2 \leq C_v B_{ijkl}(\varepsilon)(\mathbf{x}) t_{kl} t_{ij}, \quad (3.9)
\]
for all \( \varepsilon, 0 < \varepsilon \leq \varepsilon_0 \), for all \( \mathbf{x} \in \overline{\Omega} \) and all \( t = (t_{ij}) \in \mathbb{S}^2 \).

**Remark 3.4.** Note that the proof for the scaled viscosity tensor \( (B_{ijkl}(\varepsilon)) \) would follow the steps of the proof for the elasticity tensor \( (A_{ijkl}(\varepsilon)) \) in Theorem 3.3-2, [10], since from a quality point of view their expressions differ in replacing the Lamé constants by the two viscosity coefficients.
Let the scaled applied forces \( f^\varepsilon : [0, T] \times \Omega \rightarrow \mathbb{R}^3 \) and \( h^\varepsilon : [0, T] \times (\Gamma_+ \cup \Gamma_-) \rightarrow \mathbb{R}^3 \) be defined by
\[
f^\varepsilon = (f^i)(t, x) = f(\varepsilon)(t, \pi^\varepsilon(x)) \quad \forall x \in \Omega, \quad \text{where } x^\varepsilon = \pi^\varepsilon(x) \in \Omega^\varepsilon \text{ and } \forall t \in [0, T],
\]
\[
h^\varepsilon = (h^i)(t, x^\varepsilon) = h(\varepsilon)(t, x) \quad \forall x \in \Gamma_+ \cup \Gamma_-, \quad \text{where } x^\varepsilon = \pi^\varepsilon(x) \in \Gamma^\varepsilon_+ \cup \Gamma^\varepsilon_- \text{ and } \forall t \in [0, T].
\]

Also, we introduce \( u_0(\varepsilon) : \Omega \rightarrow \mathbb{R}^3 \) by
\[
u_0(\varepsilon)(x) := u_0^\varepsilon(x^\varepsilon) \quad \forall x \in \Omega, \quad \text{where } x^\varepsilon = \pi^\varepsilon(x) \in \Omega^\varepsilon,
\]
and define the space
\[
V(\Omega) := \{ v = (v_t) \in [H^1(\Omega)]^3; v = 0 \text{ on } \Gamma_0 \},
\]
which is a Hilbert space, with associated norm denoted by \( || \cdot ||_{1, \Omega} \).

We assume that the scaled applied forces are given by
\[
f(\varepsilon)(t, x) = \varepsilon^p f^p(t, x) \quad \forall x \in \Omega \text{ and } \forall t \in [0, T],
\]
\[
h(\varepsilon)(t, x) = \varepsilon^{p+1} h^{p+1}(t, x) \quad \forall x \in \Gamma_+ \cup \Gamma_- \text{ and } \forall t \in [0, T],
\]
where \( f^p \) and \( h^{p+1} \) are functions independent of \( \varepsilon \) and where \( p \) is a natural number that will show the order of the volume and surface forces, respectively. Then, the scaled variational problem can be written as follows:

**Problem 3.5.** Find \( u(\varepsilon) : [0, T] \times \Omega \rightarrow \mathbb{R}^3 \) such that,
\[
u(\varepsilon)(t, \cdot) \in V(\Omega) \quad \forall t \in [0, T],
\]
\[
\int_\Omega A^{ijkl}(\varepsilon)e_{ijkl}(\varepsilon, u(\varepsilon)(t))e_{ijkl}(\varepsilon, v)\sqrt{g(\varepsilon)}dx + \int_\Omega B^{ijkl}(\varepsilon)e_{ijkl}(\varepsilon, \dot{u}(\varepsilon)(t))e_{ijkl}(\varepsilon, v)\sqrt{g(\varepsilon)}dx
\]
\[
= \int_\Omega \varepsilon^p f^p(t) v_t \sqrt{g(\varepsilon)}dx + \int_{\Gamma_+ \cup \Gamma_-} \varepsilon^{p+1} h^{p+1}(t) v_t \sqrt{g(\varepsilon)}d\Gamma \quad \forall v \in V(\Omega), \quad \text{a.e. in } (0, T),
\]
\[
u(\varepsilon)(0, \cdot) = u_0(\varepsilon)(\cdot).
\]

From now on, for each \( \varepsilon > 0 \), we shall use the shorter notation \( e_{ijkl}(\varepsilon) \equiv e_{ijkl}(\varepsilon; u(\varepsilon)) \) and \( \dot{e}_{ijkl}(\varepsilon) \equiv \dot{e}_{ijkl}(\varepsilon; \dot{u}(\varepsilon)) \), for its time derivative. We recall the existence and uniqueness of the Problem 3.5 in the following theorem whose proof can be found in [R]:

**Theorem 3.6.** Let \( \Omega \) be a domain in \( \mathbb{R}^3 \) defined previously in this section and let \( \Theta \) be a \( C^2 \)-diffeomorphism of \( \Omega \) onto its image \( \Theta(\Omega) \), such that the three vectors \( g_i = \partial_i \Theta(x) \) are linearly independent for all \( x \in \Omega \). Let \( f^i(\varepsilon) \in L^2(0, T; L^2(\Omega)) \), \( h^i(\varepsilon) \in L^2(0, T; L^2(\Gamma_i)) \), where \( \Gamma_i := \Gamma_i^+ \cup \Gamma_i^- \). Let \( u_0(\varepsilon) \in V(\Omega) \). Then, there exists a unique solution \( u(\varepsilon) = (u_i(\varepsilon)) : [0, T] \times \Omega \rightarrow \mathbb{R}^3 \) satisfying the Problem 3.5. Moreover \( u(\varepsilon) \in W^{1,2}(0, T; V(\Omega)) \). In addition to that, if \( \dot{f}^i(\varepsilon) \in L^2(0, T; L^2(\Omega)) \), \( \dot{h}^i(\varepsilon) \in L^2(0, T; L^2(\Gamma_i)) \), then \( u(\varepsilon) \in W^{2,2}(0, T; V(\Omega)) \).
4. Technical preliminaries

Concerning geometrical and mechanical preliminaries, we shall present some theorems, which will be used in the following sections. First, we recall the Theorem 3.3-1, [10].

**Theorem 4.1.** Let \( \omega \) be a domain in \( \mathbb{R}^2 \), let \( \theta \in C^3(\bar{\omega}; \mathcal{R}^3) \) be an injective mapping such that the two vectors \( a_\alpha = \partial_\alpha \theta \) are linearly independent at all points of \( \bar{\omega} \) and let \( \varepsilon_0 > 0 \) be as in Theorem 2.3. The functions \( \Gamma^\sigma_{ij}(\varepsilon) = \Gamma^\sigma_{ij}(\varepsilon) \) and \( g(\varepsilon) \) are defined in (3.1)–(3.2), the functions \( b_{\alpha\beta}, b^\sigma_\alpha, \Gamma^\sigma_{\alpha\beta}, a, \) are defined in Section 3 and the covariant derivatives \( b^\sigma_{\beta\alpha} \) are defined by

\[
b^\sigma_{\beta\alpha} := \partial_\alpha b^\sigma_\beta + \Gamma^\sigma_{\alpha\tau} b^\tau_\beta - \Gamma^\sigma_{\alpha\beta} b^\tau_\tau.
\]

The functions \( b_{\alpha\beta}, b^\sigma_\alpha, \Gamma^\sigma_{\alpha\beta}, b^\sigma_{\beta\alpha} \) and \( a \) are identified with functions in \( C^0(\bar{\Omega}) \). Then

\[
\Gamma^\sigma_{\alpha\beta}(\varepsilon) = \Gamma^\sigma_{\alpha\beta} - \varepsilon x_3 b^\sigma_{3\alpha} + O(\varepsilon^2), \quad \Gamma^\sigma_{\alpha\beta}(\varepsilon) = b_{\alpha\beta} - \varepsilon x_3 b^\sigma_{3\alpha} b_{\beta\sigma},
\]

\[
\partial_3 \Gamma^\sigma_{\alpha\beta}(\varepsilon) = O(\varepsilon), \quad g(\varepsilon) = -b^\sigma_\alpha - \varepsilon x_3 b^\sigma_\alpha b^\sigma_\tau + O(\varepsilon^2),
\]

\[
\Gamma^\sigma_{\alpha\beta}(\varepsilon) = \Gamma^\sigma_{\alpha\beta}(\varepsilon) = 0, \quad g(\varepsilon) = a + O(\varepsilon),
\]

for all \( 0 < \varepsilon \leq \varepsilon_0 \), where the order symbols \( O(\varepsilon) \) and \( O(\varepsilon^2) \) are meant with respect to the norm \( || \cdot ||_{0,\infty,\Omega} \) defined by

\[
||w||_{0,\infty,\Omega} = \sup\{|w(x)|; x \in \bar{\Omega}\}.
\]

Finally, there exist constants \( a_0, g_0 \) and \( g_1 \) such that

\[
0 < a_0 \leq a(y) \quad \forall y \in \bar{\omega},
\]

\[
0 < g_0 \leq g(\varepsilon)(x) \leq g_1 \quad \forall x \in \bar{\Omega} \quad \forall \varepsilon, 0 < \varepsilon \leq \varepsilon_0.
\]

**Remark 4.2.** The asymptotic behaviour of \( g(\varepsilon) \) and the contravariant components of elasticity and viscosity tensors, \( A^{ijkl}(\varepsilon) \), \( B^{ijkl}(\varepsilon) \) also implies that

\[
A^{ijkl}(\varepsilon) \sqrt{g(\varepsilon)} = A^{ijkl}(0) \sqrt{a} + \varepsilon \tilde{A}^{ijkl,1} + \varepsilon^2 \tilde{A}^{ijkl,2} + o(\varepsilon^2),
\]

\[
B^{ijkl}(\varepsilon) \sqrt{g(\varepsilon)} = B^{ijkl}(0) \sqrt{a} + \varepsilon \tilde{B}^{ijkl,1} + \varepsilon^2 \tilde{B}^{ijkl,2} + o(\varepsilon^2),
\]

for certain regular contravariant components \( \tilde{A}^{ijkl,\alpha}, \tilde{B}^{ijkl,\alpha} \) of certain tensors.

We now include the following result that will be used repeatedly in what follows (see Theorem 3.4-1, [10], for details).

**Theorem 4.3.** Let \( \omega \) be a domain in \( \mathbb{R}^2 \) with boundary \( \gamma \), let \( \Omega = \omega \times (-1, 1) \), and let \( g \in L^p(\Omega), p > 1 \), be a function such that

\[
\int_{\Omega} g \partial_3 v dx = 0, \quad \text{for all } v \in C^\infty(\bar{\Omega}) \text{ with } v = 0 \text{ on } \gamma \times [-1, 1].
\]

Then \( g = 0 \).
Remark 4.4. This result holds if \( \int_{\Omega} g \partial_\gamma v \, dx = 0 \) for all \( v \in H^1(\Omega) \) such that \( v = 0 \) in \( \Gamma_0 \). It is in this way that we will use this result in the following.

We now introduce the average with respect to the transversal variable, which plays a major role in this study. To that end, let \( \mathbf{v} \) represent real or vectorial functions defined almost everywhere over \( \Omega = \omega \times (-1, 1) \). We define the transversal average as

\[
\bar{v}(\mathbf{y}) = \frac{1}{2} \int_{-1}^{1} v(\mathbf{y}, x_3) \, dx_3
\]

for almost all \( \mathbf{y} \in \omega \). Given \( \eta = (\eta_i) \in [H^1(\omega)]^3 \), let

\[
\gamma_{\alpha\beta}(\eta) := \frac{1}{2} (\partial_{\beta} \eta_\alpha + \partial_\alpha \eta_\beta) - \Gamma_{\alpha\beta}^\sigma \eta_\sigma - h_{\alpha\beta} \eta_3,
\]

(4.4)
denote the covariant components of the linearized change of metric tensor associated with a displacement field \( \mathbf{a} \) of the surface \( S \). Next theorem will show some results related with the transversal averages that will be useful in what follows.

Theorem 4.5. Let \( \omega \) be a domain in \( \mathbb{R}^2 \), let \( \Omega = \omega \times (-1, 1) \) and \( T > 0 \).

(a) Let \( v \in W^{1,2}(0, T, \mathcal{L}^2(\Omega)) \). Then \( v(\mathbf{y}) \) is finite for almost all \( \mathbf{y} \in \omega \), belongs to \( W^{1,2}(0, T, \mathcal{L}^2(\omega)) \), and

\[
||v||_{W^{1,2}(0, T, \mathcal{L}^2(\omega))} \leq \frac{1}{\sqrt{2}} ||v||_{W^{1,2}(0, T, \mathcal{L}^2(\Omega))}.
\]

If \( \partial_3 v = 0 \) in the distributions sense \( (\int_{\Omega} v \partial_3 \varphi \, dx = 0 \ \forall \varphi \in \mathcal{D}(\Omega)) \) then \( v \) does not depend on \( x_3 \) and

\[
v(\mathbf{y}, x_3) = \bar{v}(\mathbf{y}) \text{ for almost all } (\mathbf{y}, x_3) \in \Omega.
\]

(b) Let \( v \in W^{1,2}(0, T, H^1(\Omega)) \). Then \( v \in W^{1,2}(0, T, H^1(\omega)) \), \( \partial_\alpha \bar{v} = \bar{\partial_\alpha v} \) and

\[
||\bar{v}||_{W^{1,2}(0, T, H^1(\omega))} \leq \frac{1}{\sqrt{2}} ||v||_{W^{1,2}(0, T, H^1(\Omega))}.
\]

Let \( \gamma_0 \) be a subset \( \partial \gamma \)-measurable of \( \gamma \). If \( v = 0 \) on \( \gamma_0 \times [-1, 1] \) then \( \bar{v} = 0 \) on \( \gamma_0 \); in particular, \( \bar{v} \in W^{1,2}(0, T, H^1_0(\omega)) \) if \( v \equiv 0 \) on \( \gamma \times [-1, 1] \).

(c) Let \( (v(\varepsilon))_{\varepsilon > 0} \) be a sequence of functions \( v(\varepsilon) \in W^{1,2}(0, T, H^1(\Omega)) \) and let \( \bar{v} \in W^{1,2}(0, T, \mathcal{L}^2(\omega)) \) such that

\[
\partial_3 v(\varepsilon) \to 0 \text{ in } W^{1,2}(0, T, \mathcal{L}^2(\Omega)) \text{ and } \bar{v}(\varepsilon) \to \bar{v} \text{ in } W^{1,2}(0, T, \mathcal{L}^2(\omega)) \text{ when } \varepsilon \to 0.
\]

Then, \( v(\varepsilon) \to \bar{v} \) in \( W^{1,2}(0, T, \mathcal{L}^2(\Omega)) \) when \( \varepsilon \to 0 \), where the function \( \bar{v} \in W^{1,2}(0, T, \mathcal{L}^2(\omega)) \) is identified with a function in \( W^{1,2}(0, T, \mathcal{L}^2(\Omega)) \) taking \( \bar{v}(\mathbf{y}, x_3) := \bar{v}(\mathbf{y}) \) for almost all \( (\mathbf{y}, x_3) \in \Omega \).
(d) Let $\theta \in C^3(\bar{\omega}; \mathbb{R}^3)$. We consider a sequence $(v(\varepsilon))_{\varepsilon > 0}$ of vector fields $v(\varepsilon) = (v_i(\varepsilon)) \in W^{1,2}(0, T, [H^1(\Omega)]^3)$ that is bounded in $W^{1,2}(0, T, [L^2(\Omega)]^3)$. Then,

$$\left(\frac{e_{a||\beta}(\varepsilon; v(\varepsilon))}{\gamma_{a\beta}(v(\varepsilon))}\right) \to 0 \text{ in } W^{1,2}(0, T, L^2(\omega)) \text{ when } \varepsilon \to 0.$$ 

Now, we shall introduce two results that will be needed for the convergence result. Given $v = (v_i) \in [H^1(\omega)]^3$ let us define

$$\gamma_{a\beta}(v) := \frac{1}{2}(\partial_{\beta}v_{\alpha} + \partial_{\alpha}v_{\beta}) - \Gamma^\sigma_{a\alpha}v_{\sigma} - b_{a\beta}v_3,$$

$$\rho_{a\beta}(v) := \partial_{\alpha}v_3 - \Gamma^\sigma_{a\beta}\partial_{\sigma}v_3 - b^\sigma_{a\beta}v_3 + b^\sigma_{a}(\partial_{\beta}v_{\sigma} - \Gamma^\gamma_{\beta\sigma}v_{\gamma}) + b^\gamma_{3a}(\partial_{\alpha}v_{\gamma} - \Gamma^\gamma_{\alpha\gamma}v_{\gamma}) + b^\gamma_{3\beta}v_{\gamma},$$

$$e^i_{a||b}(\varepsilon; v) := \frac{1}{\varepsilon}\gamma_{a\beta}(v) + \frac{1}{\varepsilon}\partial_{\beta}e_{a||\beta}(\varepsilon; v) + \frac{1}{\varepsilon}\partial_{\alpha}e_{a||\alpha}(\varepsilon; v).$$

**Theorem 4.6.** Let the functions $\Gamma^\alpha_{a\beta}, b_{a\beta}, b^\beta_a \in C^0(\bar{\omega})$ be identified with functions in $C^0(\bar{\Omega})$ and we consider $\varepsilon_0$ defined as in Theorem 2.3. Then there exists a constant $\tilde{C} > 0$ such that for all $\varepsilon, 0 < \varepsilon \leq \varepsilon_0$ and all $v \in W^{1,2}(0, T, [H^1(\Omega)]^3)$, the scaled linearized strains $e_{a||\beta}(\varepsilon; v)$ satisfy

$$\left|\frac{1}{\varepsilon}e^i_{a||b}(\varepsilon; v) - e_{a||\beta}(\varepsilon; v)\right|_{W^{1,2}(0, T, L^2(\Omega))} \leq \tilde{C}\varepsilon \sum_{\alpha} |v_{\alpha}|_{W^{1,2}(0, T, L^2(\Omega))},$$

$$\left|\frac{1}{\varepsilon}\partial_{\beta}e_{a||\beta}(\varepsilon; v) + \rho_{a\beta}(v)\right|_{W^{1,2}(0, T, H^{-1}(\Omega))} \leq \tilde{C} \left(\sum_{i} |e^i_{a||b}(\varepsilon; v)|_{W^{1,2}(0, T, L^2(\Omega))} + \varepsilon \sum_{\alpha} |v_{\alpha}|_{W^{1,2}(0, T, L^2(\Omega))} \right).$$

**Theorem 4.7.** Let $(u(\varepsilon))_{\varepsilon > 0}$ be a sequence of functions $u(\varepsilon) \in W^{1,2}(0, T, V(\Omega))$ that satisfies

$$u(\varepsilon) \to u \text{ in } W^{1,2}(0, T, [H^1(\Omega)]^3),$$

$$\frac{1}{\varepsilon}e_{a||b}(\varepsilon; u(\varepsilon)) \to e_{a||b} \text{ in } W^{1,2}(0, T, L^2(\Omega)),$$

when $\varepsilon \to 0$. Then $u$ is independent of the transversal variable $x_3$. Moreover, $u \in H^1(\omega) \times H^2(\omega)$ with $\bar{u}_i = \partial_{\nu}u_3 = 0$ on $\gamma_0$, $\gamma_{a\beta}(u) = 0$, $\rho_{a\beta}(u) \in W^{1,2}(0, T, L^2(\Omega))$ and $\rho_{a\beta}(u) = -\partial_{b}e^i_{a||b}$. If in addition, there exist functions $\kappa_{a\beta} \in W^{1,2}(0, T, H^{-1}(\Omega))$ such that $\rho_{a\beta}(u(\varepsilon)) \to \kappa_{a\beta}$ in $W^{1,2}(0, T, H^{-1}(\Omega))$ as $\varepsilon \to 0$, then

$$u(\varepsilon) \to u \text{ in } W^{1,2}(0, T, [H^1(\Omega)]^3),$$

$$\rho_{a\beta}(u) = \kappa_{a\beta} \text{ hence, } \kappa_{a\beta} \in W^{1,2}(0, T, L^2(\Omega)).$$
Remark 4.8. The Theorems 4.3, 4.6, and 4.7 are corollaries of Theorems 4.2-1, 5.2-1 and 5.2-2, [10], respectively, and their proofs follow straightforward from the results presented there. The main difference is that we are interested in obtaining the corresponding conclusions in the Bochner spaces, namely $W^{1,2}(0, T, L^2(\Omega))$, $W^{1,2}(0, T, L^2(\omega))$, $W^{1,2}(0, T, H^1(\omega))$, $W^{1,2}(0, T, H^{-1}(\Omega))$ or $W^{1,2}(0, T, [H^1(\Omega)]^3)$. Therefore, most of the changes of the proof consists in adding an additional integral with respect to the time variable and proving the statements for the functions and their time derivatives, alternately, over the spaces $L^2(0, T, L^2(\Omega))$, $L^2(0, T, L^2(\omega))$, $L^2(0, T, H^1(\Omega))$, $L^2(0, T, H^1(\omega))$, $L^2(0, T, H^{-1}(\Omega))$ or $L^2(0, T, [H^1(\Omega)]^3)$.

Finally, in the next theorem we recall a three-dimensional inequality of Korn’s type for a family of viscoelastic shells, that can also be found in Theorem 5.3-1, [10].

Theorem 4.9. Assume that $\theta \in C^3(\bar{\omega} ; \mathbb{R}^3)$ and we consider $\varepsilon_0$ defined as in Theorem 2.3. We consider a family of viscoelastic membrane shells with thickness $2\varepsilon$ each having the same middle surface $S = \theta(\bar{\omega})$ and with each subjected to a boundary condition of place along a portion of its lateral face having the same set $\theta(\gamma_0)$ as its middle curve. Then there exist a constant $\varepsilon_1$ verifying $0 < \varepsilon_1 < \varepsilon_0$ and a constant $C > 0$ such that, for all $\varepsilon$, $0 < \varepsilon \leq \varepsilon_1$, the following three-dimensional inequality of Korn’s type holds,

$$||v||_{1, \Omega} \leq \frac{C}{\varepsilon} \left( \sum_{i,j} |c_{i,j}(\varepsilon; v)|^2_{0, \Omega} \right)^{1/2} \forall v = (v_i) \in V(\Omega). \quad (4.5)$$

5. Completion spaces and Admissible forces

In this section we shall introduce ad hoc spaces which complete the original ones presented in the obtention of the two-dimensional equations of the viscoelastic membrane shell problem in [8], leading to well posed problems. Moreover, we also shall introduce some assumptions needed on the applied forces. Recall that,

$$V_F(\omega) := \{ \eta = (\eta_i) \in H^1(\omega) \times H^1(\omega) \times H^2(\omega); \eta_i = \partial_\nu \eta_i = 0 \text{ on } \gamma_0, \gamma_{\alpha\beta}(\eta) = 0 \text{ in } \omega \}.$$

In [8], we justified the two-dimensional equations of the viscoelastic membrane shells, where the middle surface $S$ is elliptic and the boundary condition of place is considered on the whole lateral face of the shell. These assumptions lead to $V_F(\omega) = \{0\}$ (see [9] for details). Therefore in this paper, we shall considered the remaining cases where some of those assumptions are not verified but still $V_F(\omega) = \{0\}$, known as the generalized membrane shells equations. Let us define the spaces

$$V(\omega) := \{ \eta = (\eta_i) \in [H^1(\omega)]^3; \eta_i = 0 \text{ on } \gamma_0 \}, \quad (5.1)$$

$$V_0(\omega) := \{ \eta = (\eta_i) \in V(\omega); \gamma_{\alpha\beta}(\eta) = 0 \text{ in } \omega \}, \quad (5.2)$$

$$V_K(\omega) := \{ \eta = (\eta_i) \in H^1(\omega) \times H^1(\omega) \times H^2(\omega); \eta_i = \partial_\nu \eta_i = 0 \text{ on } \gamma_0 \}. \quad (5.3)$$
and also, we introduce the seminorms defined by

$$|v|^M|_{\Omega} := (|\partial_\beta v|^2 + (|v|^{M}|_{\Omega})^2)^{1/2} \forall v \in V(\Omega),$$

$$|\eta|^M|_{\omega} := \left(\sum_{\alpha,\beta} |\gamma_{\alpha\beta}(\eta)|^2_{\alpha,\beta}\right)^{1/2} \forall \eta = (\eta_i) \in H^1(\omega) \times H^1(\omega) \times L^2(\omega).$$

Since $V_F(\omega) = \{0\}$ by assumption, the seminorm $| \cdot |^M$ is a norm over the space $V_K(\omega)$.

Now, we shall distinguish two different subsets of generalized membrane shells, depending on whether or not the space $V_0(\omega)$ contains only the zero function. One of the difficulties faced is the introduction of abstract spaces, which do not have any physical meaning. We consider a generalized membrane shell of first kind when $V_0(\omega) = \{0\}$ (hence, $V_F(\omega) = \{0\}$), this is, when the seminorm $| \cdot |^M$ is a norm over the space $V(\omega)$ (hence, will be a norm over $V_K(\omega) \subset V(\omega)$). Therefore, the abstract spaces are defined by

$$V^# M(\Omega) := \text{completion of } V(\Omega) \text{ with respect to } | |^M, \quad (5.4)$$

$$V^# M(\omega) := \text{completion of } V(\omega) \text{ with respect to } | |^M. \quad (5.5)$$

**Remark 5.1.** Let consider the space,

$$V_0(\Omega) := \{v \in [H^1(\Omega)]^3; v = 0 \text{ on } \Gamma_0, \partial_\beta v = 0 \text{ in } \Omega, \gamma_{\alpha\beta}(\bar{v}) = 0 \text{ in } \omega\}. $$

Hence, by the definition of the seminorm $| \cdot |^M$ we deduce that a generalized membrane shell is of first kind if and only if $V_0(\Omega) = \{0\}$.

Otherwise, if $V_0(\omega) \neq \{0\}$ but still $V_F(\omega) = \{0\}$, this is, if $| \cdot |^M$ is a norm over $V_K(\omega)$ but not over $V(\omega)$, the shell is a generalized membrane of second kind. Therefore, the abstract spaces are defined by

$$V^# M(\Omega) := \text{completion of } V(\Omega)/V_0(\Omega) \text{ with respect to } | |^M, \quad (5.6)$$

$$V^# M(\omega) := \text{completion of } V(\omega)/V_0(\omega) \text{ with respect to } | |^M. \quad (5.7)$$

**Remark 5.2.** Notice that, in both cases, these “abstract” spaces may not be space of distributions.

Now, we shall present which are the assumptions needed for the applied forces in this problem. Let us define for each $\varepsilon > 0$, the real function $L(\varepsilon)(t) : V(\Omega) \longrightarrow \mathbb{R}$ given by

$$L(\varepsilon)(t)(v) := \int_{\Omega} f^i(t) v_i \sqrt{g(\varepsilon)} dx + \int_{\Gamma_+ \cup \Gamma_-} h^i(t) v_i \sqrt{g(\varepsilon)} d\Gamma, \quad (5.8)$$

$\forall v \in V(\Omega), \forall t \in [0, T]$, with $f^i \in L^2(0, T; L^2(\Omega))$ and $h^i \in L^2(0, T; L^2(\Gamma_+ \cup \Gamma_-))$. It is easy to check that this function is continuous with respect the norm $|| \cdot ||_{1,\Omega}$ and uniform with
respect $0 < \varepsilon \leq \varepsilon_0$, with $\varepsilon_0$ defined in Theorem 2.3. By the inequality of Korn’s type in Theorem 4.9, there exists a constant $K$ such that

$$|L(\varepsilon)(t)(\mathbf{v})| \leq \frac{K(t)}{\varepsilon} \left( \sum_{i,j} |e_{i\varepsilon,j}(\varepsilon; \mathbf{v})|^2 \right)^{1/2},$$

for all $\mathbf{v} \in V(\Omega)$ and $t \in [0, T]$.

Therefore, $L(\varepsilon)(t)$ is also continuous with respect to the norm defined by

$$\mathbf{v} \mapsto \left( \sum_{i,j} |e_{i\varepsilon,j}(\varepsilon; \mathbf{v})|^2 \right)^{1/2},$$

but not uniform with respect to $\varepsilon$ for the applied forces is made. Notice that $V(\Omega)$ is a Hilbert space with respect the interior product,

$$(\mathbf{v}, \mathbf{w}) := \int_{\Omega} e_{i\varepsilon,j}(\varepsilon; \mathbf{v}) e_{i\varepsilon,j}(\varepsilon; \mathbf{w}) \sqrt{g(\varepsilon)} dx, \quad \forall \mathbf{v}, \mathbf{w} \in V(\Omega),$$

since it is easy to verify that the norm (5.9) satisfies the parallelogram’s equality. Then, applying the Riez’s Representation Theorem, there exists a $G(t) \in [H^1(\Omega)]^3$ for each $t \in [0, T]$, such that

$$L(\varepsilon)(t)(\mathbf{v}) = \int_{\Omega} e_{i\varepsilon,j}(\varepsilon; \mathbf{G}(t)) e_{i\varepsilon,j}(\varepsilon; \mathbf{v}) \sqrt{g(\varepsilon)} dx \forall \mathbf{v} \in V(\Omega).$$

Therefore, let us define $F^{ij}(\varepsilon)(t) := e_{i\varepsilon,j}(\varepsilon; \mathbf{G}(t))$ for each $t \in [0, T]$, such that

$$L(\varepsilon)(t)(\mathbf{v}) = \int_{\Omega} F^{ij}(\varepsilon)(t) e_{i\varepsilon,j}(\varepsilon; \mathbf{v}) \sqrt{g(\varepsilon)} dx \forall \mathbf{v} \in V(\Omega).$$

If $|F^{ij}(\varepsilon)|_{0,\Omega}$ is uniformly bounded with respect to $\varepsilon$, we ensure the uniform continuity of the linear form. Moreover, we need $F^{ij}(\varepsilon)$ to have a limit in $L^2(0, T; L^2(\Omega))$, when $\varepsilon \to 0$. Therefore, the applied forces over a family of generalized membranes will be known as admissible forces, if for each $\varepsilon > 0$ there exist functions $F^{ij}(\varepsilon) = F^{ji}(\varepsilon) \in L^2(0, T; L^2(\Omega))$ and $F^{ij} = F^{ji} \in L^2(0, T; L^2(\Omega))$ such that the equality (5.12) holds for all $\varepsilon$, $0 < \varepsilon \leq \varepsilon_0$ and $F^{ij}(\varepsilon) \to F^{ij}$ in $L^2(0, T, L^2(\Omega))$ when $\varepsilon \to 0$. Therefore, if the applied forces are admissible, there exists a constant $K_0(t) > 0$ such that,

$$|L(\varepsilon)(t)(\mathbf{v})| \leq K_0(t) \left( \sum_{i,j} |e_{i\varepsilon,j}(\varepsilon; \mathbf{v})|^2 \right)^{1/2},$$

(5.13)

Thus, we need to assume additional hypotheses for the contravariant components $f^{i\varepsilon} \in L^2(0, T; L^2(\Omega))$, $h^{i\varepsilon} \in L^2(0, T; L^2(\Gamma_+ \cup \Gamma_-))$. So that, the right-hand side of the equation (2.10) can be written for each $\varepsilon > 0$ and for all $t \in [0, T]$ as follows:

$$\int_{\Omega^\varepsilon} f^{i\varepsilon}(t) \sqrt{g} dx^\varepsilon + \int_{\Gamma^+_\varepsilon \cup \Gamma^-_\varepsilon} h^{i\varepsilon}(t) \sqrt{g} d\Gamma^\varepsilon = \varepsilon L(\varepsilon)(\mathbf{v})(t),$$

(5.14)
Remark 5.3. Notice that, by considering this expression we are making an assumption on the orders of the applied forces. Actually, these orders are the ones corresponding to the viscoelastic membrane shell equations obtained in [8], that is, \( p = 0 \) in the Problem 3.5.

Then, we can write the equations in the reference domain by taking into account the definition of the admissible forces.

Problem 5.4. Find \( u(\varepsilon) : (0, T) \times \Omega \rightarrow \mathbb{R}^3 \) such that

\[
\int_\Omega A_{ijkl}(\varepsilon) e_{k[i]}(\varepsilon, u(\varepsilon)(t)) e_{l]}(\varepsilon, v) \sqrt{g(\varepsilon)} dx + \int_\Omega B_{ijkl}(\varepsilon) e_{k[i]}(\varepsilon, \dot{u}(\varepsilon)(t)) e_{l]}(\varepsilon, v) \sqrt{g(\varepsilon)} dx = L(\varepsilon)(v) \quad \forall v \in V(\Omega), \ a.e. \ in \ (0, T),
\]

\[
\begin{align*}
&u(\varepsilon)(0, \cdot) = u_0(\varepsilon)(\cdot). \\
&u(\varepsilon) \in V(\Omega) \quad \forall t \in [0, T],
\end{align*}
\]

(5.15)

The Problem 5.4 is a particular case of the Problem 3.5 hence we can ensure the existence, uniqueness and regularity of solution announced in the Theorem 3.6, taking into account the admissible forces defined above.

6. Asymptotic Analysis as \( \varepsilon \to 0 \). Convergence results

To begin with, we reformulate the two-dimensional membrane shell problem obtained in [8] taking into account the admissible forces and the abstract spaces defined in the previous section. Let us remind the definition of the two-dimensional fourth-order tensors that appeared naturally in that study,

\[
a^{\alpha\beta\sigma\tau} := \frac{2\rho^2 + 4\mu\theta^2}{\theta + \rho} a^{\alpha\beta} a^{\sigma\tau} + 2\mu(a^{\alpha\sigma} a^{\beta\tau} + a^{\alpha\tau} a^{\beta\sigma}),
\]

(6.1)

\[
b^{\alpha\beta\sigma\tau} := \frac{2\theta\rho}{\theta + \rho} a^{\alpha\beta} a^{\sigma\tau} + \rho(a^{\alpha\sigma} a^{\beta\tau} + a^{\alpha\tau} a^{\beta\sigma}),
\]

(6.2)

\[
c^{\alpha\beta\sigma\tau} := \frac{2(\theta \Lambda)^2}{\theta + \rho} a^{\alpha\beta} a^{\sigma\tau},
\]

(6.3)

where

\[
\Lambda := \left( \frac{\lambda}{\theta} - \frac{\lambda + 2\mu}{\theta + \rho} \right).
\]

(6.4)

For the sake of brevity, we only consider viscoelastic generalized membrane shells of the first kind, as those of the second kind are treated in a similar fashion. This leads to the scaled two-dimensional variational problem of a viscoelastic generalized membrane shell of first kind:
Problem 6.1. Find $\xi(t, \cdot) \in V^\#_M(\omega) \forall t \in [0, T]$ such that,

\[
B^\#_M(\xi, \eta) = L^\#_M(\eta) \quad \forall \eta \in V^\#_M(\omega), \text{ a.e. } t \in (0, T), \\
\xi(0, \cdot) = \xi_0(\cdot), \tag{6.5}
\]

where $B^\#_M$ and $L^\#_M$ are the unique continuous extensions from $W^{1,2}(0, T; V(\omega))$ to $W^{1,2}(0, T, V^\#_M(\omega))$ and from $V(\omega)$ to $V^\#_M(\omega)$ of the functions $B_M : W^{1,2}(0, T, V(\omega)) \times V(\omega) \rightarrow \mathbb{R}$ and $L_M : V(\omega) \rightarrow \mathbb{R}$, respectively, defined by

\[
B_M(\xi, \eta) := \int_\omega a^{\alpha\beta\sigma\tau} \gamma_{\sigma\tau}(\xi(t)) \gamma_{\alpha\beta}(\eta) \sqrt{a} \, dy + \int_\omega b^{\alpha\beta\sigma\tau} \gamma_{\sigma\tau}(\xi(t)) \gamma_{\alpha\beta}(\eta) \sqrt{a} \, dy \\
- \int_0^t e^{-k(t-s)} \int_\omega c^{\alpha\beta\sigma\tau} \gamma_{\sigma\tau}(\xi(s)) \gamma_{\alpha\beta}(\eta) \sqrt{a} \, dy \, ds, \\
L_M(\eta) := \int_\omega \varphi^{\alpha\beta}(t) \gamma_{\alpha\beta}(\eta) \sqrt{a} \, dy, 	ag{6.6}
\]

where we introduced the constant $k$ defined by

\[
k := \frac{\lambda + 2\mu}{\theta + \rho}, \tag{6.8}
\]

and where $\varphi^{\alpha\beta}$ is an auxiliary function, related with the admissible forces, that will appear naturally in this study, given by

\[
\varphi^{\alpha\beta}(t) := \int_{-1}^1 \left( F^{\alpha\beta}(t) - \frac{\theta}{\theta + \rho} F^{33}(t) \alpha^{\alpha\beta} + \frac{\theta \Lambda}{\theta + \rho} \int_0^t e^{-k(t-s)} F^{33}(s) \, ds \alpha^{\alpha\beta} \right) \, dx_3, 	ag{6.9}
\]

for all $t \in [0, T]$.

The Problem 6.1 is well posed and it has existence and uniqueness of solution. The proof of the next theorem follows using the same arguments made in the proof of the existence and uniqueness of solution of the de-scaled problem of the viscoelastic membrane shell (see Theorem 4.10, [8]).

Theorem 6.2. Let $\omega$ be a domain in $\mathbb{R}^2$, let $\theta \in C^2(\bar{\omega}; \mathbb{R}^3)$ be an injective mapping such that the two vectors $a_\alpha = \partial_\alpha \theta$ are linearly independent at all points of $\bar{\omega}$. Let $\varphi^{\alpha\beta} \in L^2(0, T; L^2(\omega))$ and $\xi_0 \in V^\#_M(\omega)$. Then the Problem 6.1 has a unique solution $\xi \in W^{1,2}(0, T; V^\#_M(\omega))$. In addition to that, if $\varphi^{\alpha\beta} \in L^2(0, T; L^2(\omega))$, then $\xi \in W^{2,2}(0, T; V^\#_M(\omega))$.

For each $\varepsilon > 0$, we assume that the initial condition for the scaled linear strains is

\[
e_{ij}(\varepsilon)(0, \cdot) = 0, \tag{6.10}
\]

this is, that the domain is on its natural state with no strains on it at the beginning of the period of observation.
Now, we present here the main result of this paper, that the scaled three-dimensional unknown, \( \mathbf{u}(\varepsilon) \), converges as \( \varepsilon \) tends to zero towards a limit \( \mathbf{u} \). Moreover, its transversal average, \( \mathbf{u}(\varepsilon) \), converges as \( \varepsilon \) tends to zero to the solution \( \mathbf{\xi} = \hat{\mathbf{u}} \) of the two-dimensional Problem 5.1 posed over the set \( \omega \). Given \( \mathbf{v} \in L^2(0,T,[L^2(\Omega)]^3) \) and \( \mathbf{\eta} \in L^2(0,T,[L^2(\omega)]^3) \), we shall use the notation
\[
|\mathbf{v}|^M_{T,\Omega} := \left( \int_0^T (|\mathbf{v}(t)|^M_{\Omega})^2 \, dt \right)^{1/2} , \quad |\mathbf{\eta}|^M_{T,\omega} := \left( \int_0^T (|\mathbf{\eta}(t)|^M_{\omega})^2 \, dt \right)^{1/2} .
\]

**Theorem 6.3.** Let us suppose that \( \mathbf{\theta} \in C^3(\bar{\omega} ; \mathbb{R}^3) \) and let \( \varepsilon_0 \) be defined as in Theorem 2.3. Consider a family of generalized membrane shells of first kind with thickness \( 2\varepsilon \), having each one the same middle surface \( \mathbf{S} = \mathbf{\theta} \), under a boundary condition of place along a portion of its lateral face, with the same set \( \mathbf{\theta}(\gamma_0) \) as the middle curve and subjected to admissible forces. Let \( \mathbf{u}(\varepsilon) \) be for every \( \varepsilon \), \( 0 < \varepsilon \leq \varepsilon_0 \) the solution of the three-dimensional problem under admissible forces in Problem 5.1. Then there exists \( \mathbf{u} \in W^{1,2}(0,T, V^\#_M(\omega)) \) and \( \mathbf{\xi} \in W^{1,2}(0,T, V^\#_M(\omega)) \) such that \( \mathbf{u}(\varepsilon) \to \mathbf{u} \) in \( W^{1,2}(0,T, V^\#_M(\omega)) \) when \( \varepsilon \to 0 \). Moreover,
\[
\mathbf{\xi}(\varepsilon) := \frac{1}{2} \int_{-1}^1 \mathbf{u}(\varepsilon) dx_3 \to \mathbf{\xi} \quad \text{in} \quad W^{1,2}(0,T, V^\#_M(\omega)) \quad \text{when} \quad \varepsilon \to 0 .
\]

Furthermore, the limit \( \mathbf{\xi} \) satisfies the Problem 6.1.

**Proof.** We follow the same structure of the proof of the Theorem 5.6-1, [10], so we will reference those steps which apply in the same manner.

(i) There exists \( \varepsilon_2 \), \( 0 < \varepsilon_2 \leq \varepsilon_0 \) and a constant \( c_0 \) such that, for all \( 0 < \varepsilon \leq \varepsilon_2 \),
\[
|\mathbf{v}|^M_{\Omega} \leq c_0 \left( \sum_{i,j} |e_{ij}^j(\varepsilon; \mathbf{v})|^2_{0,\Omega} \right)^{1/2} \quad \forall \mathbf{v} \in V(\Omega) .
\]

The proof can be found in step (i) in Theorem 5.6-1, [10], so we omit it.

(ii) *A priori boundedness and extractions of weakly convergent sequences.* The seminorms \( |\mathbf{u}(\varepsilon)|^M_{T,\Omega} \) and \( |\mathbf{u}(\varepsilon)|^M_{T,\omega} \), the seminorms of the respective time derivatives and the norms \( ||\mathbf{u}(\varepsilon)||_{W^{1,2}(0,T,[H^1(\Omega)]^3)} \) and \( |e_{ij}^j(\varepsilon)|_{W^{1,2}(0,T,L^2(\Omega))} \) are bounded independently of \( \varepsilon \), \( 0 < \varepsilon \leq \varepsilon_2 \).

Furthermore, by the definition of the spaces \( V^\#_M(\Omega) \) and \( V^\#_M(\omega) \) in (5.4)–(5.5), there exists a subsequence, also denoted by \( (\mathbf{u}(\varepsilon))_{\varepsilon > 0} \), and there exist \( \mathbf{u} \in W^{1,2}(0,T, V^\#_M(\Omega)) \), \( \mathbf{u}^{-1} = (\mathbf{u}^{-1}) \in W^{1,2}(0,T, V(\Omega)) \), \( e_{ij}^j \in W^{1,2}(0,T, L^2(\Omega)) \) and \( \mathbf{\xi} \in W^{1,2}(0,T, V^\#_M(\omega)) \) such that
\[
\mathbf{u}(\varepsilon) \to \mathbf{u} \quad \text{in} \quad W^{1,2}(0,T, V^\#_M(\Omega)) ,
\]
\[
\varepsilon \mathbf{u}(\varepsilon) \to \mathbf{u}^{-1} \quad \text{in} \quad W^{1,2}(0,T, [H^1(\Omega)]^3) \quad \text{hence}, \quad \varepsilon \mathbf{u}(\varepsilon) \to \mathbf{u}^{-1} \quad \text{in} \quad W^{1,2}(0,T, [L^2(\Omega)]^3) ,
\]
\[
e_{ij}^j(\varepsilon) \to e_{ij}^j \quad \text{in} \quad W^{1,2}(0,T, L^2(\Omega)) ,
\]
\[
\partial_t u_3(\varepsilon) = \varepsilon e_{3ij}(\varepsilon) \to 0 \quad \text{in} \quad W^{1,2}(0,T, L^2(\Omega)) ,
\]
\[
\mathbf{u}(\varepsilon) \to \mathbf{\xi} \quad \text{in} \quad W^{1,2}(0,T, V^\#_M(\omega)) ,
\]
when \( \varepsilon \rightarrow 0 \).

Let \( v = u(\varepsilon) \) in (5.15) and we find that
\[
\int_\Omega A^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx + \frac{1}{2} \frac{\partial}{\partial t} \int_\Omega B^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx = L_t(\varepsilon)(u(\varepsilon)),
\]
a.e. \( t \in (0, T) \). Integrating over the interval \([0, T]\), using (3.9) and (6.10) we obtain that
\[
\int_0^T \left( \int_\Omega A^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx \right) dt \leq \int_0^T L(\varepsilon)(u(\varepsilon))dt. \tag{6.13}
\]
Now, by (5.13) and the Cauchy-Schwartz inequality
\[
\int_0^T L(\varepsilon)(u(\varepsilon))dt \leq K_0 \int_0^T \left( \sum_{i,j} |e_{i\|j}(\varepsilon)|^2_{0,\Omega} \right)^{1/2} dt
\]
\[
\leq K_0 \sqrt{T} \left( \int_0^T \left( \sum_{i,j} |e_{i\|j}(\varepsilon)|^2_{0,\Omega} \right) dt \right)^{1/2} \tag{6.14}
\]
On the other hand, by (3.11), (3.8) and step (i)
\[
c_0^{-2} C_0^{-1} g_0^{-1/2} \left( |u(\varepsilon)|^2_{T,\Omega} \right) \leq C_0^{-1} g_0^{-1/2} \int_0^T \left( \sum_{i,j} |e_{i\|j}(\varepsilon)|^2_{0,\Omega} \right) dt
\]
\[
\leq \int_0^T \left( \int_\Omega A^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx \right) dt. \tag{6.15}
\]
Now, (6.13)–(6.15) together imply that \( |e_{i\|j}(\varepsilon)|_{L^2(0,T;L^2(\Omega))} \) is bounded and, as a consequence, \( |u(\varepsilon)|^M_{T,\Omega} \) and \( |\dot{u}(\varepsilon)|^M_{T,\Omega} \) do as well. By the Theorem 4.9 it follows that \( ||u(\varepsilon)||_{L^2(0,T;H^1(\Omega)^3)} \) is bounded.

Let \( v = \dot{u}(\varepsilon) \) in (5.15) and we find that
\[
\frac{1}{2} \frac{\partial}{\partial t} \int_\Omega A^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx + \int_\Omega B^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx = L(\varepsilon)(\dot{u}(\varepsilon)),
\]
a.e. \( t \in (0, T) \). Integrating over the interval \([0, T]\), using (3.8) and (6.10) we obtain that
\[
\int_0^T \left( \int_\Omega B^{ijkl}(\varepsilon)e_{k\|l}(\varepsilon)e_{i\|j}(\varepsilon)\sqrt{g(\varepsilon)}dx \right) dt \leq \int_0^T L (\varepsilon)(\dot{u}(\varepsilon))dt, \tag{6.16}
\]
that is analogous to (6.13) with the contravariant components of the viscosity tensor instead. Hence, using similar arguments and (3.9), we find that \( |\dot{e}_{i\|j}(\varepsilon)|_{L^2(0,T;L^2(\Omega))} \) are bounded and, as a consequence, \( |\dot{u}(\varepsilon)|^M_{T,\Omega} \) and \( |\dot{u}(\varepsilon)|^M_{T,\Omega} \) do as well. By the Theorem 4.9 it follows that \( ||\dot{u}(\varepsilon)||_{L^2(0,T;H^1(\Omega)^3)} \) is bounded. Therefore, the a priori boundedness and convergences announced in this step are verified.
(iii) We obtain the expressions for the limits $e_{i\|j}$ found in the previous step. Let $v = (v_i) \in V(\Omega)$. Then, by the definitions (5.31)-(5.37),

$$\varepsilon e_{\alpha\|\beta}(\varepsilon; v) \to 0 \text{ in } L^2(\Omega),$$

$$\varepsilon e_{\alpha\|3}(\varepsilon; v) \to \frac{1}{2} \partial_3 v_\alpha \text{ in } L^2(\Omega),$$

$$\varepsilon e_{3\|3}(\varepsilon; v) = \partial_3 v_3 \text{ for all } \varepsilon > 0.$$  

Let $v = \varepsilon v \in V(\Omega)$ in (5.15), we find that

$$\int_\Omega A_{ijkl}^{ij}(\varepsilon) \left( \varepsilon e_{k\|l}(\varepsilon) e_{i\|j}(\varepsilon, v) \right) \sqrt{g(\varepsilon)} \, dx + \int_\Omega B_{ijkl}^{ij}(\varepsilon) \left( \varepsilon \dot{e}_{k\|l}(\varepsilon) e_{i\|j}(\varepsilon, v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$= \int_\Omega \left( A^{\alpha\beta\gamma\delta}(\varepsilon) e_{\sigma\|\tau}(\varepsilon) + A^{\alpha\beta\gamma\delta}(\varepsilon) e_{3\|3}(\varepsilon) \right) \left( \varepsilon e_{\alpha\|\beta}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$+ \int_\Omega \left( 4A^{\alpha\beta\gamma\delta}(\varepsilon) e_{\sigma\|3}(\varepsilon) \right) \left( \varepsilon e_{\alpha\|3}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$+ \int_\Omega \left( A^{\alpha\beta\gamma\delta}(\varepsilon) e_{\sigma\|\tau}(\varepsilon) + A^{\alpha\beta\gamma\delta}(\varepsilon) e_{3\|3}(\varepsilon) \right) \left( \varepsilon e_{3\|3}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$+ \int_\Omega \left( B^{\alpha\beta\gamma\delta}(\varepsilon) e_{\sigma\|\tau}(\varepsilon) + B^{\alpha\beta\gamma\delta}(\varepsilon) e_{3\|3}(\varepsilon) \right) \left( \varepsilon e_{\alpha\|\beta}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$+ \int_\Omega \left( 4B^{\alpha\beta\gamma\delta}(\varepsilon) e_{\sigma\|3}(\varepsilon) \right) \left( \varepsilon e_{\alpha\|3}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$+ \int_\Omega \left( B^{\alpha\beta\gamma\delta}(\varepsilon) e_{\sigma\|\tau}(\varepsilon) + B^{\alpha\beta\gamma\delta}(\varepsilon) e_{3\|3}(\varepsilon) \right) \left( \varepsilon e_{3\|3}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$= \int_\Omega F^{\alpha\beta}(\varepsilon) \left( \varepsilon e_{\alpha\|\beta}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx + \int_\Omega 2F^{\alpha\beta}(\varepsilon) \left( \varepsilon e_{\alpha\|3}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx$$

$$+ \int_\Omega F^{3\|3}(\varepsilon) \left( \varepsilon e_{3\|3}(\varepsilon; v) \right) \sqrt{g(\varepsilon)} \, dx, \text{ a.e. } t \in (0, T).$$

Let $\varepsilon \to 0$. As a consequence of the asymptotic behaviour of the functions $\varepsilon e_{i\|j}(\varepsilon; v)$ above, the function $g(\varepsilon)$ and the contravariant components of the fourth order tensors $A_{ijkl}^{ij}(\varepsilon)$ and $B_{ijkl}^{ij}(\varepsilon)$ (see Theorems 4.1 and 3.3 respectively) and the weak convergences found in (ii), we obtain that

$$\int_\Omega 2\mu a^{\alpha\beta} e_{\alpha\|3}\partial_3 v_\beta + (\lambda + 2\mu) e_{3\|3}\partial_3 v_3 \sqrt{\det g(\varepsilon)} \, dx + \int_\Omega \lambda a^{\alpha\beta} e_{\alpha\|\beta}\partial_3 v_3 \sqrt{\det g(\varepsilon)} \, dx$$

$$+ \int_\Omega \rho a^{\alpha\beta} \dot{e}_{\alpha\|3}\partial_3 v_\beta + (\theta + \rho) \dot{e}_{3\|3}\partial_3 v_3 \sqrt{\det g(\varepsilon)} \, dx + \int_\Omega \theta a^{\alpha\beta} \dot{e}_{\alpha\|\beta}\partial_3 v_3 \sqrt{\det g(\varepsilon)} \, dx$$

$$= \int_\Omega \left( F^{3\|3}\partial_3 v_3 + F^{33}\partial_3 v_3 \right) \sqrt{\det g(\varepsilon)} \, dx, \text{ a.e. } t \in (0, T).$$

(6.17)

Let $v \in V(\Omega)$ be independent of $x_3$. Then, we have that

$$\int_\Omega 2\mu a^{\alpha\beta} e_{\alpha\|3}\partial_3 v_\beta \sqrt{\det g(\varepsilon)} \, dx + \int_\Omega \rho a^{\alpha\beta} \dot{e}_{\alpha\|3}\partial_3 v_\beta \sqrt{\det g(\varepsilon)} \, dx = \int_\Omega \left( F^{3\|3}\partial_3 v_3 \right) \sqrt{\det g(\varepsilon)} \, dx.$$

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Hence, by Theorem 4.3 this equation leads to,

\[ 2\mu a^{\alpha\sigma} e_{\alpha|3} + \rho a^{\alpha\sigma} \dot{e}_{\alpha|3} = F^{\alpha 3}, \]

and using that \((a_{\alpha\sigma})^{-1} = a^{\alpha\sigma}\), we obtain the following ordinary differential equation,

\[ 2\mu e_{\alpha|3} + \rho \dot{e}_{\alpha|3} = a^{\alpha\sigma} F^{\sigma 3}. \]  

(6.18)

**Remark 6.4.** Note that removing time dependency and viscosity, that is taking \(\rho = 0\), the equation leads to the one studied in [10], that is, the elastic case.

In order to solve the equation (6.18) in the more general case, we assume that the viscosity coefficient \(\rho\) is strictly positive. Moreover, we can prove that this equation is equivalent to

\[ \frac{\partial}{\partial t} \left( e^{\frac{2\mu t}{\rho}} e_{\alpha|3}(t) \right) = \frac{1}{\rho} a_{\alpha\sigma} e^{\frac{2\mu t}{\rho}} F^{\sigma 3}(t). \]

Integrating with respect to the time variable and using (6.10) we find that

\[ e_{\alpha|3}(t) = \frac{1}{\rho} a_{\alpha\sigma} \int_0^t e^{-\frac{2\mu}{\rho}(t-s)} F^{\sigma 3}(s) ds \text{ in } \Omega, \forall t \in [0, T]. \]  

(6.19)

Moreover, from (6.18) we obtain that,

\[ \dot{e}_{\alpha|3}(t) = \frac{1}{\rho} \left( a_{\alpha\sigma} F^{\sigma 3}(t) - 2\mu e_{\alpha|3}(t) \right) \text{ in } \Omega, \text{ a.e. } t \in (0, T). \]

(6.20)

Now, taking \(v \in V(\Omega)\) such that \(v_{\alpha} = 0\), then we have that

\[
\int_\Omega (\lambda + 2\mu) e_{3|3} \partial_3 v_3 \sqrt{\alpha} dx + \int_\Omega \lambda a^{\alpha\beta} e_{\alpha|\beta|3} \partial_3 v_3 \sqrt{\alpha} dx \\
+ \int_\Omega (\theta + \rho) \dot{e}_{3|3} \partial_3 v_3 \sqrt{\alpha} dx + \int_\Omega \theta a^{\alpha\beta} \dot{e}_{\alpha|\beta|3} \partial_3 v_3 \sqrt{\alpha} dx = \int_\Omega F^{33} \partial_3 v_3 \sqrt{\alpha} dx.
\]

(6.20)

Applying Theorem 4.3 we obtain the following differential equation,

\[ \lambda^{\alpha\beta} e_{\alpha|\beta} + (\lambda + 2\mu) e_{3|3} + \theta a^{\alpha\beta} \dot{e}_{\alpha|\beta} + (\theta + \rho) \dot{e}_{3|3} = F^{33}. \]  

(6.21)

**Remark 6.5.** Once again, note that removing time dependency and viscosity, that is taking \(\theta = \rho = 0\), the equation leads to the one studied in [10], that is, the elastic case.

In order to solve the equation (6.21) in the more general case, we assume that the viscosity coefficient \(\theta\) is strictly positive. Moreover, we can prove that this equation is equivalent to

\[ \theta e^{-\frac{\lambda t}{\rho}} \frac{\partial}{\partial t} \left( a^{\alpha\beta} e_{\alpha|\beta}(t) e^{\frac{\lambda t}{\rho}} \right) = F^{33}(t) - (\theta + \rho) e^{-\frac{\lambda t+2\mu}{\rho t}} \frac{\partial}{\partial t} \left( e_{3|3}(t) e^{\frac{\lambda t+2\mu}{\rho t}} \right). \]  

(6.22)
Integrating respect to the time variable and using (6.10) we find that,
\[
e^{3 |3| (t)} = \frac{1}{\theta + \rho} \int_0^t e^{\frac{3|3|}{\theta + \rho} s} F^{33}(s) ds - \frac{\theta}{\theta + \rho} \int_0^t e^{(\frac{3|3|}{\theta + \rho} - \frac{\lambda}{\theta + \rho}) s} \partial_s \left( a^{\alpha\beta} e_{\alpha|\beta}(s) e^{\frac{3|3|}{\theta + \rho} s} \right) ds,
\]
now integrating by parts and simplifying we conclude that,
\[
e^{3 |3| (t)} = \frac{1}{\theta + \rho} \int_0^t e^{-k(t-s)} F^{33}(s) ds - \frac{\theta}{\theta + \rho} \left( a^{\alpha\beta} e_{\alpha|\beta}(t) + \Lambda \int_0^t e^{-k(t-s)} a^{\alpha\beta} e_{\alpha|\beta}(s) ds \right)
\]
in $\Omega$ , $\forall \ t \in [0, T]$, and where $\Lambda$ and $k$ are defined in (6.4) and (6.8), respectively. Moreover, from (6.21) we obtain that,
\[
\dot{e}^{3|3| (t)} = \frac{1}{\theta + \rho} F^{33}(t) - \frac{\lambda}{\theta + \rho} a^{\alpha\beta} e_{\alpha|\beta}(t) - \frac{\lambda + 2\mu}{\theta + \rho} e^{3|3| (t)} - \frac{\theta}{\theta + \rho} a^{\alpha\beta} \dot{e}_{\alpha|\beta}(t),
\]
in $\Omega$ , a.e. $t \in (0, T)$.

(iv) The family $\left( u(\varepsilon) \right)_{\varepsilon > 0}$ verifies
\[
\left( e^{\alpha|\beta}(\varepsilon) - \gamma_{\alpha\beta}(u(\varepsilon)) \right) \rightarrow 0 \ \text{in} \ W^{1,2}(0, T, L^2(\omega)) \ \text{when} \ \varepsilon \rightarrow 0. \tag{6.23}
\]
As a consequence, the subsequence considered in (ii) verifies
\[
\gamma_{\alpha\beta}(u(\varepsilon)) \rightarrow \overline{e^{\alpha|\beta}} \ \text{in} \ W^{1,2}(0, T, L^2(\omega)). \tag{6.24}
\]
This proof is a corollary of the step (iv) in Theorem 5.6-1, [10]. We follow the same arguments made there and use Theorem 4.5 (a) and (b), to prove that those convergences, and the corresponding ones of the time derivatives, hold in the space $L^2(0, T, L^2(\omega))$. Then, the conclusion follows.

(v) We obtain the equations satisfied by the limits $e_{\alpha|\beta}$ found in the step (ii).
Let $\mathbf{v} = (v_i) \in V(\Omega)$ be independent of the transversal variable $x_3$. Then, by the definitions (3.5)–(3.7),
\[
e_{\alpha|\beta}(\varepsilon; \mathbf{v}) \rightarrow \gamma_{\alpha\beta}(\mathbf{v}) \ \text{in} \ L^2(\Omega),
\]
\[
e_{\alpha|\beta}(\varepsilon; \mathbf{v}) \rightarrow \frac{1}{2} \partial_\alpha v_3 + b_\alpha^\sigma v_\sigma \ \text{in} \ L^2(\Omega),
\]
\[
e_{3|3}(\varepsilon; \mathbf{v}) = 0 \ \text{for all} \ \varepsilon > 0.
\]
Keep such a function $\mathbf{v} \in V(\Omega)$ in (5.15). Taking the limit when $\varepsilon \rightarrow 0$ in the right-hand side of the equation (5.15), we have that
\[
\lim_{\varepsilon \rightarrow 0} L(\varepsilon)(\mathbf{v}) = \int_\Omega \left( F^{\alpha\beta} \gamma_{\alpha\beta}(\mathbf{v}) + 2F^{\alpha3} \left( \frac{1}{2} \partial_\alpha v_3 + b_\alpha^\sigma v_\sigma \right) \right) \sqrt{a} dx. \tag{6.25}
\]
Also, let $\varepsilon \to 0$ in the left-hand side of the equation (6.15). By the asymptotic behaviour of functions $g(\varepsilon)$ and the contravariant components of the fourth order tensors $A^{ijkl}(\varepsilon)$ and $B^{ijkl}(\varepsilon)$ (see Theorem 4.1 and 3.3 respectively), the convergences of the strain components $e_{iij}(\varepsilon; v)$ above and the weak convergences of $e_{iij}(\varepsilon) \rightharpoonup e_{iij}$ in $W^{1,2}(0, T, L^2(\Omega))$ from (ii), we find that,

$$
\int_{\Omega} A^{ijkl}(0) e_{iij}(\varepsilon; v) \sqrt{a} dx + \int_{\Omega} B^{ijkl}(0) e_{iij}(\varepsilon; v) \sqrt{a} dx
$$

which, using the relations found in (iii) yields that,

$$
\int_{\Omega} \left( \left( \lambda - \frac{\theta}{\theta + \rho} (\theta A + \lambda) \right) a^{\alpha \beta} a^{\sigma \tau} + \mu (a^{\alpha \sigma} a^{\beta \tau} + a^{\alpha \tau} a^{\beta \sigma}) \right) e_{\sigma \tau \gamma \alpha \beta}(\varepsilon; v) \sqrt{a} dx
$$

+ $\int_{\Omega} \frac{\theta \lambda}{\theta + \rho} \int_{0}^{t} e^{-k(t-s)} F^{33}(s) ds a^{\alpha \beta} a^{\sigma \tau} \gamma_{\alpha \beta}(\varepsilon; v) \sqrt{a} dx
$$

+ $\int_{\Omega} \left( \frac{\theta}{\theta + \rho} a^{\alpha \beta} a^{\sigma \tau} + \frac{\rho}{2} (a^{\alpha \sigma} a^{\beta \tau} + a^{\alpha \tau} a^{\beta \sigma}) \right) e_{\sigma \tau \gamma \alpha \beta}(\varepsilon; v) \sqrt{a} dx
$$

+ $\int_{\Omega} \left( \frac{\theta}{\theta + \rho} F^{33} a^{\alpha \beta} \gamma_{\alpha \beta}(\varepsilon; v) \sqrt{a} dx - \int_{\Omega} \frac{(\theta A)^2}{\theta + \rho} \int_{0}^{t} e^{-k(t-s)} a^{\sigma \tau} e_{\sigma \tau \gamma \alpha \beta}(\varepsilon; v) \sqrt{a} dx
$$

+ $\int_{\Omega} 2 F^{33} \left( \frac{1}{2} \partial_{a} v_{3} + b_{\alpha} v_{\sigma} \right) \sqrt{a} dx,
$$

(6.26)

Then, simplifying this expression is equivalent to

$$
\frac{1}{2} \int_{\Omega} a^{\alpha \beta \sigma \tau} e_{\sigma \tau \gamma \alpha \beta}(\varepsilon; v) \sqrt{a} dx + \frac{1}{2} \int_{\Omega} b^{\alpha \beta \sigma \tau} e_{\sigma \tau \gamma \alpha \beta}(\varepsilon; v) \sqrt{a} dx
$$

- $\frac{1}{2} \int_{0}^{t} e^{-k(t-s)} \int_{\Omega} c^{\alpha \beta \sigma \tau} e_{\sigma \tau \gamma \alpha \beta}(s) \sqrt{a} dx ds
$$

+ $\int_{\Omega} \frac{\theta \lambda}{\theta + \rho} \int_{0}^{t} e^{-k(t-s)} F^{33}(s) ds a^{\alpha \beta} a^{\sigma \tau} \gamma_{\alpha \beta}(\varepsilon; v) \sqrt{a} dx
$$

+ $\int_{\Omega} \frac{\theta}{\theta + \rho} F^{33} a^{\alpha \beta} \gamma_{\alpha \beta}(\varepsilon; v) \sqrt{a} dx + \int_{\Omega} 2 F^{33} \left( \frac{1}{2} \partial_{a} v_{3} + b_{\alpha} v_{\sigma} \right) \sqrt{a} dx,
$$

where $a^{\alpha \beta \sigma \tau}$, $b^{\alpha \beta \sigma \tau}$ and $c^{\alpha \beta \sigma \tau}$ denote the contravariant components of the two-dimensional...
fourth order tensors defined in (6.1)–(6.3). Hence, together with (6.25) leads to
\[
\int_\omega a^{\alpha\beta\gamma\delta} v^{\alpha|\tau} v^{\beta|\sigma} \sqrt{\omega} dy + \int_\omega b^{\alpha\beta\gamma\delta} v^{\alpha|\tau} v^{\beta|\sigma} \sqrt{\omega} dy \\
- \int_0^t e^{-k(t-s)} \int_\omega a^{\alpha\beta\gamma\delta} v^{\alpha|\tau} (s) v^{\beta|\sigma} \sqrt{\omega} dy ds \\
= \int_\omega \int_{-1}^1 F^{\alpha\beta} dx_3 v^{\alpha\beta}(\bar{v}) \sqrt{\omega} dy - \int_\omega \int_{-1}^1 \frac{\theta \Lambda}{\theta + \rho} \int_0^t e^{-k(t-s)} F^{33}(s) ds a^{\alpha\beta} dx_3 v^{\alpha\beta}(\bar{v}) \sqrt{\omega} dy \\
- \int_\omega \int_{-1}^1 \frac{\theta}{\theta + \rho} F^{33} a^{\alpha\beta} dx_3 v^{\alpha\beta}(\bar{v}) \sqrt{\omega} dy = \int_\omega \varphi^{\alpha\beta} v^{\alpha\beta}(\bar{v}) \sqrt{\omega} dy.
\] (6.27)
where \(\varphi^{\alpha\beta}\) denotes the real function defined in (6.9). Now, given \(\eta \in V(\omega)\), there exists a function \(v \in V(\Omega)\) independent of \(x_3\) such that \(v = \eta\). Hence the announced relations hold for all \(\eta \in V(\omega)\), a.e. \(t \in (0, T)\).

(vi) The subsequence \((u(\varepsilon))_{\varepsilon>0}\) from (ii) satisfies
\[
\varepsilon u(\varepsilon) \rightharpoonup 0 \text{ in } W^{1,2}(0, T, [H^1(\Omega)]^3),
\]
\[
\partial_3 u_{\alpha}(\varepsilon) \rightharpoonup 0 \text{ in } W^{1,2}(0, T, L^2(\Omega)),
\]
when \(\varepsilon \to 0\). Moreover, \(e_{\alpha||\beta}\) are independent of the transversal variable \(x_3\).

This proof is a corollary of the step (vi) in Theorem 5.6-1, [10]. By the step (ii) the functions \(u^{-1}(\varepsilon) := \varepsilon u(\varepsilon) \in W^{1,2}(0, T, V(\Omega))\) satisfy
\[
\frac{1}{\varepsilon} e_{\alpha||\beta}(\varepsilon; u^{-1}(\varepsilon)) \rightharpoonup e_{\alpha||\beta} \text{ in } W^{1,2}(0, T, L^2(\Omega)).
\]
Hence, by Theorem 4.7, \(u^{-1} \in V_F(\omega)\) and consequently \(u^{-1} = 0\), since \(V_F(\omega) = \{0\}\) by assumption. By the same result, \(u^{-1}\) is independent of \(x_3\), hence \(u^{-1} = 0\) and the first weak convergence is proved. Moreover, this implies that \(\varepsilon u(\varepsilon) \rightharpoonup 0 \text{ in } W^{1,2}(0, T, [L^2(\Omega)]^3)\). Now, by (5.6) we have that
\[
\partial_3 u_{\alpha} = 2\varepsilon e_{\alpha||\beta}(\varepsilon) - \varepsilon \partial_3 u_{3\beta}(\varepsilon) + 2\varepsilon \Gamma_{\alpha\beta}(\varepsilon) u_\sigma(\varepsilon).
\]
Therefore, together with the convergences in (ii) and above and the boundedness of the sequence \((\Gamma_{\alpha\beta}(\varepsilon))_{\varepsilon>0}\) in \(C^0(\Omega)\) by the Theorem 2.8 imply that \(\partial_3 u_{\alpha}(\varepsilon) \rightharpoonup 0 \text{ in } W^{1,2}(0, T, L^2(\Omega))\). Moreover, since \(u^{-1} = 0\) and (6.30)–(6.31), taking \(u = u^{-1}\) in Theorem 4.7 we have that \(\partial_3 e_{\alpha||\beta} = -\rho_{\alpha\beta}(u^{-1}) = 0\). Therefore, the functions \(e_{\alpha||\beta}\) are independent of \(x_3\).

(vii) The following strong convergences are satisfied,
\[
e_{\alpha||\beta} \rightharpoonup e_{\alpha||\beta} \text{ in } W^{1,2}(0, T, L^2(\Omega)),
\]
\[
\varepsilon u(\varepsilon) \rightharpoonup 0 \text{ in } W^{1,2}(0, T, [H^1(\Omega)]^3),
\]
\[
\gamma_{\alpha\beta}(\bar{u}(\varepsilon)) \rightharpoonup \gamma_{\alpha\beta}(\bar{u}) \text{ in } W^{1,2}(0, T, L^2(\omega)),
\]
\[
\frac{\bar{u}(\varepsilon)}{\varepsilon} \rightharpoonup \xi \text{ in } W^{1,2}(0, T, V_{\xi}^*(\omega)),
\]
for all $t \in [0, T]$.

Let us define,

$$\Psi(\varepsilon) := \int_{\Omega} A^{ijkl}(\varepsilon)(e_{kli}(\varepsilon) - e_{kli}(\varepsilon)) \varepsilon d\tau$$

$$+ \int_{\Omega} B^{ijkl}(\varepsilon)(\dot{e}_{kli}(\varepsilon) - \dot{e}_{kli}(\varepsilon)) \varepsilon d\tau$$

$$= L(\varepsilon)(u(\varepsilon)) - \int_{\Omega} A^{ijkl}(\varepsilon)(2e_{kli}(\varepsilon) - e_{kli})e_{iij} \varepsilon d\tau$$

$$+ \int_{\Omega} B^{ijkl}(\varepsilon)((\dot{e}_{kli}e_{iij} - \frac{\partial}{\partial t}(e_{kli}(\varepsilon)e_{iij})) \varepsilon d\tau, \text{ a.e. } t \in (0, T).$$

We have that,

$$\int_{\Omega} A^{ijkl}(\varepsilon)(e_{kli}(\varepsilon) - e_{kli})(e_{iij}(\varepsilon) - e_{iij}) \varepsilon d\tau$$

$$+ \frac{1}{2} \frac{\partial}{\partial t} \int_{\Omega} B^{ijkl}(\varepsilon)(e_{kli}(\varepsilon) - e_{kli})(e_{iij}(\varepsilon) - e_{iij}) \varepsilon d\tau = \Psi(\varepsilon), \text{ a.e. } t \in (0, T).$$

Integrating over the interval $[0, T]$, using (3.9) and (6.10) we find that

$$\int_{0}^{T} \left( \int_{\Omega} A^{ijkl}(\varepsilon)(e_{kli}(\varepsilon) - e_{kli})(e_{iij}(\varepsilon) - e_{iij}) \varepsilon d\tau \right) dt \leq \int_{0}^{T} \Psi(\varepsilon)dt. \quad (6.32)$$

Now, by (4.1) and (3.8)

$$C_{\varepsilon}^{-1} g_{\alpha}^{1/2} \sum_{i,j} |e_{iij}(\varepsilon) - e_{iij}|_{0,\Omega}^{2} \leq \int_{\Omega} A^{ijkl}(\varepsilon)(e_{kli}(\varepsilon) - e_{kli})(e_{iij}(\varepsilon) - e_{iij}) \varepsilon d\tau.$$

Therefore, together with the previous inequality leads to

$$C_{\varepsilon}^{-1} g_{\alpha}^{1/2} \int_{0}^{T} \left( \sum_{i,j} |e_{iij}(\varepsilon) - e_{iij}|_{0,\Omega}^{2} \right) dt \leq \int_{0}^{T} \Psi(\varepsilon)dt.$$

On the other hand, the strong convergences $F^{ij}(\varepsilon) \rightarrow F^{ij}$ in $L^{2}(0, T, L^{2}(\Omega))$ by assumption and the weak convergences $e_{iij}(\varepsilon) \rightharpoonup e_{iij}$ in $W^{1,2}(0, T, L^{2}(\Omega))$ from (ii) imply that

$$\lim_{\varepsilon \rightarrow 0} L(\varepsilon)(u(\varepsilon)) = \lim_{\varepsilon \rightarrow 0} \left( \int_{\Omega} F^{\alpha\beta}(\varepsilon)e_{\alpha\beta}(\varepsilon) \varepsilon d\tau \right)$$

$$+ \int_{\Omega} (2F^{\alpha\beta}(\varepsilon)e_{\alpha\beta}(\varepsilon) + F^{33}(\varepsilon)e_{33}(\varepsilon)) \varepsilon d\tau$$

$$= \int_{\Omega} F^{\alpha\beta}e_{\alpha\beta} \varepsilon d\tau + \int_{\Omega} (2F^{\alpha\beta}e_{\alpha\beta} + F^{33}e_{33}) \varepsilon d\tau. \quad (6.33)$$

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Now, by the asymptotic behaviour of functions $g(\varepsilon)$ and the contravariant components of the fourth order tensors $A^{ijkl}(\varepsilon)$ and $B^{ijkl}(\varepsilon)$ (see Theorem 4.1 and 3.3 respectively), we find that

$$\lim_{\varepsilon \to 0} \left( \int_{\Omega} A^{ijkl}(\varepsilon)(2e_{k||l}(\varepsilon) - e_{k||l})e_{i||j} \sqrt{g(\varepsilon)}dx \right)$$

$$- \int_{\Omega} B^{ijkl}(\varepsilon)(e_{k||l}e_{i||j} - \frac{\partial}{\partial t}(e_{k||l}(\varepsilon)e_{i||j})) \sqrt{g(\varepsilon)}dx$$

$$= \int_{\Omega} A^{ijkl}(0)e_{k||l}e_{i||j} \sqrt{g}dx + \int_{\Omega} B^{ijkl}(0)e_{k||l}e_{i||j} \sqrt{g}dx$$

$$= \int_{\Omega} (\lambda a^{\alpha\beta} a^{\sigma\tau} + \mu(a^{\alpha\sigma} a^{\beta\tau} + a^{\alpha\tau} a^{\beta\sigma})) e_{\sigma||\tau} e_{\alpha||\beta} \sqrt{g}dx + \int_{\Omega} \lambda a^{\alpha\beta} e_{3||3} e_{\alpha||\beta} \sqrt{g}dx$$

$$+ \int_{\Omega} 4\mu a^{\alpha\sigma} e_{\sigma||3} e_{\alpha||3} \sqrt{g}dx + \int_{\Omega} (\lambda a^{\sigma\tau} e_{\sigma||\tau} + (\lambda + 2\mu) e_{3||3} e_{\sigma||\tau}) \sqrt{g}dx$$

$$+ \int_{\Omega} (\theta a^{\alpha\beta} a^{\sigma\tau} + \rho(a^{\alpha\sigma} a^{\beta\tau} + a^{\alpha\tau} a^{\beta\sigma})) e_{\sigma||\tau} e_{\alpha||\beta} \sqrt{g}dx + \int_{\Omega} \theta a^{\alpha\beta} e_{3||3} e_{\alpha||\beta} \sqrt{g}dx$$

$$+ \int_{\Omega} 2\rho a^{\alpha\sigma} e_{\sigma||3} e_{\alpha||3} \sqrt{g}dx + \int_{\Omega} (\theta a^{\sigma||\tau} + (\theta + \rho) e_{3||3} e_{\sigma||\tau}) e_{3||3} \sqrt{g}dx,$$

which substituting the findings in the step (iii) and simplifying leads to

$$\int_{\Omega} a^{\alpha\beta\sigma\tau} e_{\alpha||\beta} e_{\sigma||\tau} \sqrt{g}dx + \int_{\Omega} b^{\alpha\beta\sigma\tau} e_{\alpha||\beta} e_{\sigma||\tau} \sqrt{g}dx - \int_{0}^{t} e^{-k(t-s)} \int_{\Omega} e^{\alpha\beta\sigma\tau} e_{\alpha||\beta} e_{\sigma||\tau} \sqrt{g}dx ds$$

$$+ \int_{\Omega} \int_{-1}^{1} \left( \frac{\theta}{\theta + \rho} F^{33} a^{\alpha\beta} + \frac{\theta \Lambda}{\theta + \rho} \int_{0}^{t} e^{-k(t-s)} F^{33}(s) ds a^{\alpha\beta} \right) dx e_{\alpha||\beta} \sqrt{g}$$

$$+ \int_{\Omega} (2F^{33} e_{3||3} + F^{33} e_{3||3}) \sqrt{g}dx,$$  

(6.34)

where $a^{\alpha\beta\sigma\tau}$, $b^{\alpha\beta\sigma\tau}$ and $c^{\alpha\beta\sigma\tau}$ denote the contravariant components of the fourth order tensors defined in (6.1)–(6.3). Hence, together with (6.34), we have that

$$\Psi := \lim_{\varepsilon \to 0} \Psi(\varepsilon) = \int_{\omega} \varphi^{\alpha\beta} e_{\alpha||\beta} \sqrt{g}dx - \int_{\omega} a^{\alpha\beta\sigma\tau} e_{\alpha||\beta} e_{\sigma||\tau} \sqrt{g}dx$$

$$\int_{\omega} b^{\alpha\beta\sigma\tau} e_{\alpha||\beta} e_{\sigma||\tau} \sqrt{g}dx,$$

(6.35)

with $\varphi^{\alpha\beta}$ defined in (6.9). Now, since $\bar{u}(\varepsilon) \in V(\omega)$, for each $\varepsilon > 0$, we take $\bar{v} = \eta = \bar{u}(\varepsilon)$ in (6.27) and we have that

$$\int_{\omega} a^{\alpha\beta\sigma\tau} e_{\alpha||\beta}(\bar{u}(\varepsilon)) \sqrt{g}dx + \int_{\omega} b^{\alpha\beta\sigma\tau} e_{\alpha||\beta}(\bar{u}(\varepsilon)) \sqrt{g}dx$$

$$- \int_{0}^{t} e^{-k(t-s)} \int_{\omega} c^{\alpha\beta\sigma\tau} e_{\alpha||\beta}(\bar{u}(\varepsilon)) \sqrt{g}dx ds = \int_{\omega} \varphi^{\alpha\beta} e_{\alpha||\beta}(\bar{u}(\varepsilon)) \sqrt{g}dx.$$  

(6.36)
Taking in (6.36) the limit when \( \varepsilon \to 0 \) together with the weak convergences in \((iv)\), we conclude from (6.35) that \( \Psi = 0 \). As a consequence, using the Lebesgue dominated convergence theorem in (6.32), the strong convergences \( e_{i\|j}(\varepsilon) \to e_{i\|j} \) in \( L^2(0,T,L^2(\Omega)) \) are verified. Analogously, if we define

\[
\tilde{\Psi}(\varepsilon) := \int_\Omega A^{ijkl}(\varepsilon)(e_{k\|l}(\varepsilon) - e_{k\|l})(\dot{e}_{i\|j}(\varepsilon) - \dot{\varepsilon}_{i\|j})\sqrt{g(\varepsilon)}dx
+ \int_\Omega B^{ijkl}(\varepsilon)(\dot{e}_{k\|l}(\varepsilon) - \ddot{e}_{k\|l})(\dot{e}_{i\|j}(\varepsilon) - \ddot{\varepsilon}_{i\|j})\sqrt{g(\varepsilon)}dx
= L_t(\varepsilon)(\dot{\mathbf{u}}(\varepsilon)) + \int_\Omega A^{ijkl}(\varepsilon)(e_{k\|l}\dot{e}_{i\|j} - \frac{\partial}{\partial t}(e_{k\|l}(\varepsilon)e_{i\|j}))\sqrt{g(\varepsilon)}dx
- \int_\Omega B^{ijkl}(\varepsilon)(2\dot{e}_{k\|l}(\varepsilon) - \dot{e}_{k\|l})\dot{e}_{i\|j}\sqrt{g(\varepsilon)}dx.
\]

We have that,

\[
\frac{1}{2}\frac{\partial}{\partial t} \int_\Omega A^{ijkl}(\varepsilon)(e_{k\|l}(\varepsilon) - e_{k\|l})(\dot{e}_{i\|j}(\varepsilon) - \dot{\varepsilon}_{i\|j})\sqrt{g(\varepsilon)}dx
+ \int_\Omega B^{ijkl}(\varepsilon)(\dot{e}_{k\|l}(\varepsilon) - \ddot{e}_{k\|l})(\dot{e}_{i\|j}(\varepsilon) - \ddot{\varepsilon}_{i\|j})\sqrt{g(\varepsilon)}dx = \tilde{\Psi}(\varepsilon), \text{ a.e. } t \in (0,T).
\]

Integrating over the interval \([0,T]\), using (3.8) and (6.10) we find that

\[
\int_0^T \left( \int_\Omega B^{ijkl}(\varepsilon)(\dot{e}_{k\|l}(\varepsilon) - \ddot{e}_{k\|l})(\dot{e}_{i\|j}(\varepsilon) - \ddot{\varepsilon}_{i\|j})\sqrt{g(\varepsilon)}dx \right) dt \leq \int_0^T \tilde{\Psi}(\varepsilon)dt,
\]

Now, by (3.2) and (4.1)

\[
C_v^{-1}g_0^{1/2}\sum_{i,j} \|\dot{e}_{i\|j}(\varepsilon) - \dot{\varepsilon}_{i\|j}\|^2_{0,\Omega} \leq \int_\Omega B^{ijkl}(\varepsilon)(\dot{e}_{k\|l}(\varepsilon) - \ddot{e}_{k\|l})(\dot{e}_{i\|j}(\varepsilon) - \ddot{\varepsilon}_{i\|j})\sqrt{g(\varepsilon)}dx
\]

Therefore, together with the previous inequality leads to

\[
C_v^{-1}g_0^{1/2}\int_0^T \left( \sum_{i,j} \|\dot{e}_{i\|j}(\varepsilon)(t) - \dot{\varepsilon}_{i\|j}(t)\|^2_{0,\Omega} \right) dt \leq \int_0^T \tilde{\Psi}(\varepsilon)dt,
\]

which is similar with (6.32). Therefore, using analogous arguments as before, we find that

\[
\hat{\Psi} := \lim_{\varepsilon \to 0} \tilde{\Psi}(\varepsilon) = \int_\Omega \alpha^{\alpha\beta\gamma\delta}e^{\alpha\beta\gamma\delta}_{\alpha\beta\gamma\delta}(\varepsilon)e_{\alpha\beta\gamma\delta}(\varepsilon)\sqrt{g(\varepsilon)}dx
+ \int_0^T e^{-k(t-s)} \int_\Omega \alpha^{\alpha\beta\gamma\delta}e^{\alpha\beta\gamma\delta}_{\alpha\beta\gamma\delta}(s)e_{\alpha\beta\gamma\delta}(t)\sqrt{g(\varepsilon)}dy \, ds, \text{ a.e. } t \in (0,T).
\]
Therefore, we conclude that

\[
\text{Taking in (6.39) the limit when } \varepsilon \to 0 \text{ together with the weak convergences in (iv), we conclude from (6.38) that } \hat{\Psi} = 0. \text{ As a consequence, using the Lebesgue dominated convergence theorem in (6.37), the strong convergences } \hat{\varepsilon}_{ij}(\varepsilon) \rightarrow \hat{\varepsilon}_{ij} \text{ in } L^2(0, T, L^2(\Omega)) \text{ are satisfied. Therefore, we conclude that } e_{ij}(\varepsilon) \rightarrow e_{ij} \text{ in } W^{1,2}(0, T, L^2(\Omega)).
\]

Now, let \( v = u^{-1}(\varepsilon) = \varepsilon u(\varepsilon) \). Then, using the second inequality in Theorem 4.6 we find by the step (ii) that there exist two constants \( \hat{C}, \hat{C} > 0 \) such that

\[
||\partial_3 e_{a||b}(\varepsilon) + \rho_{a\beta}(u^{-1}(\varepsilon))||_{W^{1,2}(0,T,H^{-1}(\Omega))} \leq \hat{C} \left( \varepsilon \sum_i |e_{ij}(\varepsilon)||_{W^{1,2}(0,T,L^2(\Omega))} + \varepsilon ||\varepsilon u_{3}(\varepsilon)||_{W^{1,2}(0,T,H^1(\Omega))} \right) \leq \hat{C} \varepsilon.
\]

Moreover, since \( e_{a||b}(\varepsilon) \rightarrow e_{a||b} \) in \( W^{1,2}(0, T, L^2(\Omega)) \) and the functions \( e_{a||b} \) are independent of \( x_3 \) (see step (iv)) we have that

\[
\partial_3 e_{a||b}(\varepsilon) \rightarrow \partial_3 e_{a||b} = 0 \text{ in } W^{1,2}(0, T, H^{-1}(\Omega)),
\]

hence, from the previous inequality,

\[
\rho_{a\beta}(u^{-1}(\varepsilon)) \rightarrow 0 \text{ in } W^{1,2}(0, T, H^{-1}(\Omega)).
\]

Now, applying Theorem 4.7 we have that

\[
u^{-1}(\varepsilon) = \varepsilon u(\varepsilon) \rightarrow 0 \text{ in } W^{1,2}(0, T, [H^1(\Omega)]^3).
\]

By the Theorem 4.5 (a), the strong convergences \( e_{ij}(\varepsilon) \rightarrow e_{ij} \) in \( W^{1,2}(0, T, L^2(\Omega)) \) imply that \( e_{a||b}(\varepsilon) \rightarrow e_{a||b} \) in \( W^{1,2}(0, T, L^2(\omega)) \). Therefore, by (iv)

\[
\gamma_{a\beta}(u(\varepsilon)) \rightarrow \overline{e_{a||b}} \text{ in } W^{1,2}(0, T, L^2(\omega)).
\]

As a consequence, \( \gamma_{a\beta}(u(\varepsilon))_{\varepsilon > 0} \) is a Cauchy sequence in \( W^{1,2}(0, T, L^2(\omega)) \). Now, since

\[
|u(\varepsilon) - u(\varepsilon')|^M_{\frac{1}{2}, \omega} = \left( \sum_{a,\beta} \int_0^T |\gamma_{a\beta}(u(\varepsilon))(t) - \gamma_{a\beta}(u(\varepsilon'))(t)|^M_{\omega} dt \right)^{1/2}, \text{ with } \varepsilon, \varepsilon' > 0,
\]

and the corresponding identity for the time derivatives hold, the strong convergence \( \gamma_{a\beta}(u(\varepsilon)) \rightarrow \xi \) in \( W^{1,2}(0, T, V^M_{\omega}(\omega)) \) is verified.
(viii) The limit \( \xi(t) \in V^\#_M(\omega) \) \( \forall t \in [0, T] \) found in (vii) satisfies the system of equations

\[
B^\#_M(\xi, \eta) = L^\#_M(\eta) \quad \forall \eta \in V^\#_M(\omega), \text{ a.e. } t \in (0, T),
\]

\( \xi(0, \cdot) = \xi_0(\cdot), \)

which has a unique solution. As a consequence, the convergence \( \overline{u(\varepsilon)} \to \xi \) in \( W^{1,2}(0, T, V^\#_M(\omega)) \) is verified by the all family \( (\overline{u(\varepsilon)})_{\varepsilon>0} \).

Let \( \eta \in V(\omega) \). By the steps (v) and (vii) and since \( \overline{u(\varepsilon)} \in W^{1,2}(0, T, V_M(\omega)) \), we find that

\[
\lim_{\varepsilon \to 0} B_M(\overline{u(\varepsilon)}, \eta) = \lim_{\varepsilon \to 0} \left( \int_\omega a^{\alpha\beta\sigma\tau} \gamma_{\sigma\tau}(\overline{u(\varepsilon)(t)}) \gamma_{\alpha\beta}(\eta) \sqrt{\text{ad}y} + \int_\omega b^{\alpha\beta\sigma\tau} \gamma_{\sigma\tau}(\overline{u(\varepsilon)(t)}) \gamma_{\alpha\beta}(\eta) \sqrt{\text{ad}y} \right)
\]

\[
- \int_0^t e^{-k(t-s)} \int_\omega c^{\alpha\beta\sigma\tau} \gamma_{\sigma\tau}(\overline{u(\varepsilon)(s)}) \gamma_{\alpha\beta}(\eta) \sqrt{\text{ad}y} \text{d}s
\]

\[
= \int_\omega a^{\alpha\beta\sigma\tau} e_{\sigma\tau}(t) \gamma_{\alpha\beta}(\eta) \sqrt{\text{ad}y} + \int_\omega b^{\alpha\beta\sigma\tau} e_{\sigma\tau}(t) \gamma_{\alpha\beta}(\eta) \sqrt{\text{ad}y} \]

\[
- \int_0^t e^{-k(t-s)} \int_\omega c^{\alpha\beta\sigma\tau} e_{\sigma\tau}(s) \gamma_{\alpha\beta}(\eta) \sqrt{\text{ad}y} \text{d}s = L_M(\eta).
\]

Furthermore, again by (vii) we have that

\[
\lim_{\varepsilon \to 0} B_M(\overline{u(\varepsilon)}, \eta) = B^\#_M(\xi, \eta) = L_M(\eta), \forall \eta \in V(\omega), \text{ a.e. } t \in (0, T),
\]

hence, \( B^\#_M(\xi, \eta) = L^\#_M(\eta) \) \( \forall \eta \in V^\#_M(\omega) \), a.e. \( t \in (0, T) \), by the definition of the continuous extensions \( B^\#_M \) and \( L^\#_M \). In our previous work, the reader can find the details that ensure the existence, uniqueness and regularity of the solution of this problem (see Theorem 4.10, [8]).

(ix) Let \( \Omega \) be a domain in \( \mathbb{R}^3 \). Given \( \nu = (\nu_i) \in W^{1,2}(0, T, [L^2(\Omega)]^3) \), we define the distributions,

\[
e_{ij}(\nu) := \frac{1}{2}(\partial_i \nu_j + \partial_j \nu_i) \in W^{1,2}(0, T, H^{-1}(\Omega)). \tag{6.40}
\]

Considering a sequence of functions \( \nu^k = (\nu^k) \in W^{1,2}(0, T, [L^2(\Omega)]^3) \) such that \( \nu^k \to 0 \) in \( W^{1,2}(0, T, [H^{-1}(\Omega)]^3) \) and \( e_{ij}(\nu^k) \to 0 \) in \( W^{1,2}(0, T, H^{-1}(\Omega)) \) when \( k \to \infty \). Then, \( \nu^k \to 0 \) in \( W^{1,2}(0, T, [L^2(\Omega)]^3) \).

This proof is a corollary of the step (ix) in Theorem 5.6-1, [10]. We follow the same arguments made there to prove that \( \nu^k \to 0 \) in \( L^2(0, T, [L^2(\Omega)]^3) \) and the corresponding convergences of the time derivatives in the same space. Then the conclusion follows.

(x) The following convergences are satisfied,

\[
u(\varepsilon) \to \nu \text{ in } W^{1,2}(0, T, V_M^\#(\Omega)),
\]

\[
\partial_3 u_\alpha(\varepsilon) \to 0 \text{ in } W^{1,2}(0, T, L^2(\Omega)).
\]
We shall follow similar arguments of the proof of the step (x) in Theorem 5.6-1. For the first convergence is enough to prove that \( (u(\varepsilon))_{\varepsilon > 0} \) and its time derivative, are Cauchy sequences with respect to the norm \( | \cdot |_{T,\Omega} \). By its definition we have,

\[
\int_0^T \left( |u(\varepsilon)(t) - u'(\varepsilon)(t)|_{H_0}^2 \right) dt = \int_0^T \left( \sum_{\alpha,\beta} |\gamma_{\alpha\beta}(u(\varepsilon)(t)) - \gamma_{\alpha\beta}(u'(\varepsilon)(t))|^2_{0,\omega} \right) dt \\
+ \int_0^T \left( \sum_i |\partial_3 u_i(\varepsilon)(t) - \partial_3 u_i'(\varepsilon)(t)|_{0,\omega}^2 \right) dt.
\]

and the analogous equality for the time derivative family. Since \( \partial_3 u_3(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \) and \( \gamma_{\alpha\beta}(u(\varepsilon)) \to \gamma_{\alpha\beta}(u) \) in \( W^{1,2}(0, T, L^2(\omega)) \) by steps (ii) and (vii), it only remains to prove that \( \partial_3 u_3(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \). These convergence are verified if \( \partial_3 u'(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \), with

\[
u'(\varepsilon) = (u_1(\varepsilon), u_2(\varepsilon), 0).
\]

By the step (ix), proving this is the same to showing that \( \partial_3 u'(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, [H^{-1}(\Omega)]^3) \) and \( e_{ij}(\partial_3 u'(\varepsilon)) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \) by steps (vi) and (vii), respectively, together imply that

\[
\partial_3 u_3(\varepsilon) \to 0 \text{ in } W^{1,2}(0, T, H^{-1}(\Omega)).
\]

Then \( \partial_3 u'(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, [H^{-1}(\Omega)]^3) \) holds.

For the second convergence, we first have that \( e_{33}(\partial_3 u'(\varepsilon)) = \partial_3 e_{33}(u'(\varepsilon)) = 0 \) by (6.41) and (6.42). Now, the asymptotic behaviour of the functions \( \Gamma_{\alpha\beta}^3(\varepsilon) \) (see Theorem 4.1) and the convergences \( \varepsilon e_{11}(\varepsilon) \to 0, \varepsilon u_1(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \) by steps (vi) and (vii), respectively, imply that

\[
\partial_3 u_3(\varepsilon) \to 0 \text{ in } W^{1,2}(0, T, L^2(\Omega)),
\]

hence,

\[
(\partial_3 u_3(\varepsilon) + \varepsilon \partial_3 u_3(\varepsilon) + 2\varepsilon b_\alpha' u_\sigma(\varepsilon)) \to 0 \text{ in } W^{1,2}(0, T, L^2(\Omega)),
\]

Now, since \( \partial_3 u_3(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \) and \( \varepsilon u_\sigma(\varepsilon) \to 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \) (steps (ii) and (vi), respectively) we have that

\[
2e_{\alpha\beta}(\partial_3 u'(\varepsilon)) = \partial_3 u_3(\varepsilon) \to 0 \text{ in } W^{1,2}(0, T, H^{-1}(\Omega)).
\]
Finally, by step (vi) \( e_{\alpha|\beta}(\varepsilon) \rightarrow e_{\alpha|\beta} \) in \( W^{1,2}(0, T, L^2(\Omega)) \) and \( \partial_3 e_{\alpha|\beta} = 0 \), hence we infer that \( \partial_3 e_{\alpha|\beta}(\varepsilon) \rightarrow 0 \) in \( W^{1,2}(0, T, H^{-1}(\Omega)) \). Therefore,

\[
\partial_3 e_{\alpha|\beta}(\varepsilon) = (\partial_3 e_{\alpha|\beta}(u(\varepsilon)) - \partial_3 (\Gamma^p_{\alpha\beta}(\varepsilon)u_p(\varepsilon))) \rightarrow 0 \text{ in } W^{1,2}(0, T, H^{-1}(\Omega)).
\]

Since \( \Gamma^p_{\alpha\beta}(\varepsilon) \in C^1(\bar{\Omega}) \), then \( \Gamma^p_{\alpha\beta}(\varepsilon)u_p(\varepsilon) \in W^{1,2}(0, T, H^1(\Omega)) \). Furthermore,

\[
\partial_3 (\Gamma^p_{\alpha\beta}(\varepsilon)u_p(\varepsilon)) = \partial_3 \Gamma^p_{\alpha\beta}(\varepsilon)u_p(\varepsilon) + \Gamma^p_{\alpha\beta}(\varepsilon)\partial_3 u_p(\varepsilon),
\]

\[
\Gamma^p_{\alpha\beta}(\varepsilon)\partial_3 u_p(\varepsilon) \rightarrow 0 \text{ in } W^{1,2}(0, T, H^{-1}(\Omega)),
\]

since \( \partial_3 u_p(\varepsilon) \rightarrow 0 \) in \( W^{1,2}(0, T, H^{-1}(\Omega)) \) (step (ii) and (6.43)). Now, the estimates \( ||\partial_3 \Gamma^p_{\alpha\beta}(\varepsilon)||_{0,\infty,\bar{\Omega}} \leq C\varepsilon \), with a constant \( C > 0 \) (see Theorem 6.6) and the convergences \( \varepsilon u(\varepsilon) \rightarrow 0 \) in \( W^{1,2}(0, T, [L^2(\Omega)]^3) \) (see step (vi)) imply that

\[
\partial_3 \Gamma^p_{\alpha\beta}(\varepsilon)u_p(\varepsilon) \rightarrow 0 \text{ in } W^{1,2}(0, T, L^2(\Omega)).
\]

Therefore, \( e_{\alpha\beta}(\partial_3 u'(\varepsilon)) = \partial_3 e_{\alpha\beta}(u(\varepsilon)) \rightarrow 0 \) in \( W^{1,2}(0, T, H^{-1}(\Omega)) \).

(xi) All family \( (u(\varepsilon))_{\varepsilon>0} \) converges strongly to \( u \) in the space \( W^{1,2}(0, T, V^\#_M(\Omega)) \).

The family \( (\bar{u}(\varepsilon))_{\varepsilon>0} \) converges strongly in \( W^{1,2}(0, T, V^\#_M(\omega)) \) by step (viii) and \( \partial_3 u(\varepsilon) \rightarrow 0 \) in \( W^{1,2}(0, T, L^2(\Omega)) \) for a subsequence (see steps (ii) and (x)). Then, since the limit of such sequence is unique, the whole family \( (\partial_3 u(\varepsilon))_{\varepsilon>0} \) converges in \( W^{1,2}(0, T, L^2(\Omega)) \). Therefore, by the definition of the norm \( |.|_{\partial^3 \Omega} \), \( (\bar{u}(\varepsilon))_{\varepsilon>0} \) is a Cauchy sequence in the Hilbert space \( W^{1,2}(0, T, V^\#_M(\Omega)) \), hence, the conclusion follows.

Therefore, the proof of the theorem is complete.

\[ \square \]

It remains to prove the analogous result to the previous theorem in terms of de-scaled unknowns. Therefore we need to de-scale the unknown \( \xi ; \) solution of the two-dimensional variational scaled problem. By the scaling proposed in Section 3 we define for each \( \varepsilon > 0 \) the vector field \( \xi^\varepsilon \) such that

\[
\xi^\varepsilon := \xi \text{ in } W^{1,2}(0, T, V^\#_M(\omega)).
\]

Notice that, for the viscoelastic generalized membrane shells, we can not consider the de-scaling of each component separately, since the previous equality must be understood only in the abstract completion space. Therefore, we can prove the following convergence result:

**Theorem 6.6.** Assume that \( \theta \in C^3(\tilde{\omega}; \mathbb{R}^3) \). Consider a family of viscoelastic generalized membrane shells of the first kind with thickness \( 2\varepsilon \) approaching zero and with each having the same middle surface \( S = \theta(\tilde{\omega}) \), with each subjected to a boundary condition of place along a portion of its lateral face having the same set \( \theta(\gamma_0) \) as its middle curve and subjected to admissible forces (see Section 5).
Let \( u^\varepsilon = (u_i^\varepsilon) \in W^{1,2}(0, T, V(\Omega^\varepsilon)) \) and \( \xi^\varepsilon = (\xi_i^\varepsilon) \in W^{1,2}(0, T, V_M^\#(\omega)) \) respectively denote for each \( \varepsilon > 0 \) the solutions to the three-dimensional and two-dimensional Problems 2.4 and 1.1. Moreover, let \( \xi = (\xi_i) \in W^{1,2}(0, T, V_M(\omega)) \) denote the solution to the Problem 6.1. Then we have that

\[
\xi^\varepsilon = \xi \quad \text{and} \quad \frac{1}{2\varepsilon} \int_{-\varepsilon}^{\varepsilon} u^\varepsilon dx_3^\varepsilon \rightarrow \xi \quad \text{in} \quad W^{1,2}(0, T, V_M^\#(\omega)) \quad \text{as} \quad \varepsilon \rightarrow 0,
\]

Proof. Notice that,

\[
\frac{1}{2\varepsilon} \int_{-\varepsilon}^{\varepsilon} u^\varepsilon dx_3^\varepsilon = \frac{1}{2} \int_{-1}^{1} u(\varepsilon)dx_3 = \overline{u(\varepsilon)}.
\]

Hence, the conclusion follows by applying Theorem 6.3.

7. Conclusions

We have found and mathematically justified a model for viscoelastic generalized membrane shells subjected to admissible forces. To this end we used the asymptotic expansion method (presented in our previous work [8]) and we have justified this approach by obtaining convergence theorems. As in the elastic case we have distinguished two cases (generalized membrane of the first kind or second time) depending on whether or not the space \( V_0^M(\omega) \) contains non-zero functions. For each case, completion spaces were needed in order to obtain well posed problems.

The main novelty that this model presented is a long-term memory, represented by an integral on the time variable, more specifically

\[
M(t, \eta) = \int_0^t e^{-k(t-s)} \int_\omega c_{\alpha\beta\sigma\tau}(\xi(s))\gamma_{\alpha\beta}(\eta)\sqrt{adyds},
\]

for all \( \eta \in V_M^\#(\omega) \) (analogously for second type generalized membranes). Analogous behaviour has been presented in beam models for the bending-stretching of viscoelastic rods [30], obtained by using asymptotic methods as well. Also, this kind of viscoelasticity has been described in [15, 28], for example.

As the viscoelastic case differs from the elastic case on time dependent constitutive law and external forces, we must consider the possibility that these models and the convergence result generalize the elastic case (studied in [10]). However, analogously to the asymptotic analysis made in [8], the reader can easily check, when the ordinary differential equation (6.18) and (6.21) were presented, we had to use assumptions that make it impossible to include the elastic case. Hence, the viscoelastic and elastic problems must be treated separately in order to reach reasonable and justified conclusions.

Furthermore, as in the elastic case [10], we found that \( \partial_3 u(\varepsilon)(t, \cdot) \rightarrow 0 \quad \forall \quad t \in [0, T] \), while for the elliptic case we proved that the three-dimensional limit \( u \) was independent of \( x_3 \) (see [8]). Therefore, the displacements on a viscoelastic generalized membrane shell might not be independent of the transversal variable. Moreover, notice that we proved convergence
theorems when applied admissible forces (5.12) are considered, hence, the body and surface forces can not be arbitrary as in the elliptic case.

These models together with the elliptic case presented in our previous paper [6, 7], complete the study for the viscoelastic membrane shells. As commented in this paper, it remains to study the case when $V_F(\omega)$ contains non-zero functions, whose model was presented in [8], known as the problem of a viscoelastic flexural shell. This case will be presented in a forthcoming paper [6].

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