Correlating structure with non-linear optical properties in $x\text{As}_{40}\text{Se}_{60}(1 - x)\text{As}_{40}\text{S}_{60}$ glasses

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A series of $x\text{As}_{40}\text{Se}_{60}(100 - x)\text{As}_{40}\text{S}_{60}$ glasses, where $x = 0, 25, 33, 50, 67, 75$ and $100$ mol% $\text{As}_{40}\text{Se}_{60}$, has been studied using neutron and X-ray total scattering, Raman spectroscopy and $^{77}\text{Se}$ MAS-NMR. The results are presented with measurements of non-linear refractive indices, $n_2$, and densities. There is no evidence for the formation of homopolar bonds in these glasses, but neutron correlation functions suggest that there is a non-random distribution of sulfur and selenium atoms in sulfur-rich glasses. The average number of sulfur atoms at a distance of 3–4 Å from a selenium atom, $n_{\text{SeS}}$, deviates from a linear variation with $x$ in glasses containing <50 mol% $\text{As}_{40}\text{Se}_{60}$; $n_2$ for these glasses also varies non-linearly with $x$. Importantly, a direct comparison of $n_2$ and $n_{\text{SeS}}$ gives a linear correlation, suggesting that $n_2$ may be related to the distribution of chalcogen atoms in the glasses.

1. Introduction

The low phonon energies, and large non-linear optical responses, exhibited by chalcogenide glasses enable them to transmit, and interact with, a broad spectrum of mid-infrared (mid-IR) light. These properties, coupled with an ability to act as laser hosts via the incorporation of rare earth ions, make chalcogenide glasses ideal candidates for a range of mid-IR applications, including new bright mid-IR fibre laser sources.

The As–S–Se glass system is of particular interest because melts require only moderate cooling rates (air cooling) to form a glass and it has a wide glass forming domain. The range of glass formation is bounded by the binary glass forming ranges of $\text{As}_x\text{Se}_{100-x}$, $0–62$ at% As in $\text{As}_x\text{Se}_{100-x}$, and $0–85$ at% Se in $\text{S}_x\text{Se}_{100-x}$. As–S–Se glasses show good thermal stability, with resistance to crystallisation during fibre-drawing, and non-linear refractive indices, $n_2$, > 100<$\mu$n that of silica.3,4 The effect on glass properties caused by varying the relative amounts of As, Se and S is well documented, but less well understood. For example, reported physical property measurements of $x\text{As}_{40}\text{Se}_{60}(100 - x)\text{As}_{40}\text{S}_{60}$ glasses indicate that density,5,6 and glass transition temperature, $T_g$,3,5,6 vary linearly with $x$ yet $n_2$, important for active optical applications, does not.3,4 Understanding the relationship between composition, structure and macroscopic properties will facilitate the development of glass compositions that are designed and tailored for specific optical devices.

Previous structural studies of As–Se–S glasses have predominantly used spectroscopic techniques. However, these methods have limitations in the information they can yield. $^{77}\text{Se}$ Nuclear Magnetic Resonance (NMR) studies give averaged evidence for the local environment of selenium,7,8 but sulfur and arsenic NMR are less informative as bonding information cannot be gained directly. Numerous Raman scattering studies2,3,9–11 have identified the characteristic vibrations of As–S, As–Se, S–S and Se–Se. S–S bond vibrations oscillate at relatively high wave numbers, because both atoms in the bond are light, giving rise to a band in the Raman spectra at 440–500 cm$^{-1}$.2,3,9–11 When sulfur is bonded to the heavier arsenic or selenium atoms, the bond vibrational frequency was reported to red-shift to ~270–420 cm$^{-1}$ and when the heavier oscillators (As and Se) are directly bonded (As–As, As–Se and Se–Se), bands in the wave number range 220–270 cm$^{-1}$ are observed. Attempts have been made to relate these Raman spectral features to glass structure. However, due to the similar masses of As and Se, deconvolution of contributions from As–S and Se–S at ~350 cm$^{-1}$, or As–As, Se–Se and As–Se at ~250 cm$^{-1}$, is difficult (see ref. 2 and 10 for examples). One study of As-poor glasses, where homopolar chalcogen–chalcogen (Ch–Ch) bonds are an integral part of the glass network, indicated that there is a preference for Se–Se bonds over those of S–S or S–Se, concluding that selenium chains, and [AsS$_3$] units, are more favourable than [AsSe$_3$] units and sulfur chains.10 However, Raman spectra for pure selenium and pure sulfur showed that...
Se-Se bonds have a greater oscillator strength than S–S bonds. This results in a greater intensity of the Se peaks in the Raman spectrum, and casts doubt on the conclusions drawn in ref. 10. In conclusion, difficulties in quantitatively interpreting spectroscopic measurements, coupled with a lack of structural information beyond the direct bonds in the glass, limit the use of NMR and Raman in aiding the understanding of how chalcogen atoms are distributed in chalcogenide glasses.

In contrast to spectroscopic techniques, neutron and X-ray correlation functions, \( T(r) \), yield quantitative measurements of glass structure. The results give information about the local environments of atoms in the glass, and how those structural units link together to form a 3-D glass network. Although there have been several total scattering studies\(^ {12–14} \) of the stoichiometric glasses, As\(_{20}\)Se\(_{80} \) and As\(_{40}\)S\(_{60} \), and of binary glasses with varying cation to anion ratios, As\(_{1}\)–Se\(_{2} \) and As\(_{1}\)–S\(_{2} \), there are fewer which study the effect of mixing S and Se chalcogen anions on the network structure. Furthermore, there is not much information on how glass structure underpins macroscopic glass properties.

The stoichiometric glass forming system, \( x \)As\(_{20}\)Se\(_{80} \). (100–\( x \))As\(_{40}\)S\(_{60} \) (where \( x = 0, 25, 33, 50, 67, 75, 100 \) mol%), has been chosen for this study to minimise the number of homopolar (As–As, Ch–Ch) bonds and simplify the structure.\(^ {15} \) Along this compositional tie line, an ideal glass would be solely constructed on the same nucleus), and the distinct scattering, \( I(\mathbf{Q}) \) (the interference between scattered waves from the same nucleus), and the distinct scattering, \( I(\mathbf{Q}) \) (the interference between scattered waves from different nuclei);

\[
\frac{d\sigma}{d\Omega} = I^N(\mathbf{Q}) = I^X(\mathbf{Q}) + \mathcal{I}(\mathbf{Q}).
\]

The Fourier transform of \( \mathcal{I}(\mathbf{Q}) \) yields a real-space correlation function, \( T^N(r) \). This is a measure of the frequency of the occurrence of interatomic distances in the sample and is calculated as:

\[
T^N(r) = 4\pi\rho_0^2 \left( \sum_i c_i \bar{b}_i \right)^2 + \frac{1}{\pi} \int_0^{Q_{\text{max}}} Q I^N(Q) M(Q) \sin(qr) dq
\]

where \( M(Q) \) is a modification function, such as the Lorch function,\(^ {16} \) used to reduce termination ripples in the Fourier transform and \( \rho_0 \) is the atomic number density. \( c_i \) is the atomic fraction and \( \bar{b}_i \) is the coherent neutron scattering length (see Table 1) for element \( l \). The resolution in real-space depends on the modification function and the maximum momentum transfer, \( Q_{\text{max}} \) of the experimental data.

In the case of X-ray diffraction, the scattering data has a form factor dependence. This is removed by normalising the data using the Krogh-Moe and Norman method\(^ {17,18} \) to yield the distinct scattering \( I^X(\mathbf{Q}) \). X-rays scatter from the electron cloud of an atom, resulting in a relatively broad ‘electron–electron’ correlation function. To remove this broadening and obtain a correlation function, \( T^X(r) \), that has good resolution, a sharpened distinct scattering spectra, \( \tilde{I}^X(\mathbf{Q}) \), is required. This is obtained by dividing \( \tilde{I}^X(\mathbf{Q}) \) by \( f(\mathbf{Q})^2 \), where \( f(\mathbf{Q}) \) is the mean X-ray form factor for the sample and is related to the atomic number, \( Z \), of the elements involved.\(^ {19,20} \) This new function approximates scattering from point sources;

\[
\tilde{I}^X(\mathbf{Q}) = \frac{I^X(\mathbf{Q})}{\langle f(\mathbf{Q}) \rangle^2}
\]

The X-ray correlation function, \( T^X(r) \), obtained from a Fourier transform of \( \tilde{I}^X(\mathbf{Q}) \) is given by:

\[
T^X(r) = 4\pi\rho_0^2 r + \frac{1}{\pi} \int_0^{Q_{\text{max}}} Q I^X(Q) M(Q) \sin(qr) dq
\]

Neglecting the effect of form factors on X-ray peak shapes in \( T^X(r) \), the correlation functions, \( T^N(r) \) and \( T^X(r) \), are the weighted summations over all the pairwise combinations of elements in the sample, \( \delta_{ll'} \):

\[
T^N(r) = \sum_{ll'} c_l \bar{b}_l \delta_{ll'} \rho_l(r) \quad \text{and} \quad T^X(r) = \sum_{ll'} c_l Z_l Z_{l'} \rho_{ll'}(r)
\]

### 2. Extracting coordination numbers from \( T^N(r) \)

Each peak in the neutron correlation function, \( T^N(r) \), arising from a particular pair of elements \( l \) and \( l' \). The coordination number of the peak area, \( A_{ll'} \), and the interatomic distance, \( r_{ll'} \), given by the position of the peak and according to:

\[
n_{ll'} = -\frac{r_{ll'} A_{ll'}}{(2 - \delta_{ll'}) c_l \bar{b}_l \bar{b}_{l'}}
\]

where \( \delta_{ll'} \) is the Kronecker delta and \( c_l \bar{b}_l \bar{b}_{l'} \) is the coefficient for \( \delta_{ll'} \) (I). If the coordination number for \( n_{ll'} \) and \( n_{11} \) have been calculated using the area and position of a feature in \( T^N(r) \) with a defined distance range, then the coordination numbers for \( l \) and \( l' \) atom pairs are related by the following identity:

\[
n_{ll'} c_l = n_{1l'} c_{l'}
\]

### Table 1 The variation in scattering power of As, S and Se for neutrons and X-rays

| Element | Neutron scattering length, \( b \) | X-ray form factor, \( f(\mathbf{Q}) \) |
|---------|----------------------------------|----------------------------------|
| S       | 2.847                           | 16                               |
| As      | 6.58                            | 33                               |
| Se      | 7.97                            | 34                               |
The first peak in the neutron correlation function, \(T_n(r)\), for a glass network constructed solely of [AsCh3] (Ch = S or Se) units, is comprised of directly bonded As–Ch distances. The second peak arises from the next nearest neighbours in the glass and can comprised of a mix of up to nine X···X distances (where ··· denotes a non-bonding distance between two atoms). The various contributions can be written as an expansion of eqn (5)

\[
T_n(r) = c_{As}b_{As}^2n_{AsAs} + c_{As}b_{As}b_{S}n_{AsSe} + c_{As}b_{As}b_{S}n_{AsSS} \\
+ c_{Se}b_{Se}^2n_{SeSe} + c_{Se}b_{Se}b_{S}n_{SeAS} + c_{Se}b_{Se}b_{S}n_{SeSS} \\
+ c_{bS}b_{SS}n_{SS} + c_{bS}b_{AS}n_{AS} + c_{bS}b_{Se}n_{SeSS} \tag{8}
\]

If there is a linear change in the glass structure with composition, then a weighted sum of the two end member glasses, \(x\)As40Se60 + (100 – \(x\))As40S60, will give a good approximation for the As–Ch contribution to the first peak in the neutron correlation function, and the As–As and As–Ch correlations contributing to the second peak, for a glass with composition \(x\)As40Se60 + (100 – \(x\))As40S60. However, the weighted sum will overestimate the contributions from Se···Se and S···S correlations because mixing, i.e. the presence of Se···S distances, is neglected.

The expected difference between an experimental \(T_n(r)\) and a \(T_n(r)\) predicted by a weighted sum of As40Se60 and As40S60 can be calculated by a consideration of the Ch···Ch contribution to the second peak in the experimental data. It can be shown, using eqn (6) and (8), that the expected area for a peak comprised of all the possible Ch···Ch contributions is related to the total number of chalcogen neighbours, \(n_{ChCh}\):

\[
r_{ChCh}A_{ChCh} = c_{Se}b_{Se}^2n_{SeSe} + c_{Se}b_{Se}b_{S}n_{SeAS} + c_{Se}b_{Se}b_{S}n_{SeSS} + c_{bS}b_{SS}n_{SS} + c_{bS}b_{AS}n_{AS} + c_{bS}b_{Se}n_{SeSS} \tag{9}
\]

If it is assumed that the number of chalcogen atoms at a distance \(r + \delta r\) from an Se atom, \(n_{SeCh}\), and a S atom, \(n_{SCh}\), is the same, then

\[
n_{ChCh} = n_{SS} + n_{SeS} = n_{SeSe} + n_{SeS} \tag{10}
\]

Rearranging these equations for \(n_{SS}\) and \(n_{SeS}\), substituting the result into eqn (9) and simplifying the result using the relationship in eqn (7), the following equation can be derived:

\[
r_{ChCh}A_{ChCh} = n_{ChCh}(c_{Se}b_{Se}^2 + c_{bSe}) - c_{Se}n_{SeS}(b_{Se} - b_{S})^2 \tag{11}
\]

The first term in eqn (11) is the contribution to the Ch···Ch peak from the weighted sum of Se···Se and S···S contributions. The difference between the weighted sum, and the experimental data is given by the second term and yields a peak arising from Se···S correlations. Fitting this residual allows \(n_{SeS}\) to be calculated. If there is no preference for like (or unlike) chalcogen atoms to be in the same [AsCh3] unit, then the Se-Se coordination number depends on the fraction of chalcogens which are sulfur, so that \(n_{SeS} = (1 - (x/100))n_{ChCh}\). In As40Se60 (\(x = 100\)), \(n_{SeS} = 0\) because there are no sulfur atoms in the glass. However, for As40S60S (\(x = 0\)), all the chalcogen next nearest neighbours to the single Se atom would be sulfur and \(n_{SeS} = n_{ChCh} = n_{SS}\) in As40S60. The structure of crystalline As40S60 predicts that \(n_{SS}\) for the glass should be approximately 7.33. Deviations from a linear variation in \(n_{SeS}\) from \(\sim 7.33\) to 0 with increasing \(x\) would only occur if:

1. There is a significant difference between the experimental arsenic correlations and those predicted by the weighted sum of the end members.
2. There was a preference for like anions to cluster together.
3. The number of S atoms around Se differs from the number of Se atoms around S in the distance range of interest.

### 3. Experimental setup

#### 3.1. Sample preparation

A series of glasses, \(x\)As40Se60(100 – \(x\)) As40S60, where \(x = 0, 25, 33, 50, 67, 75\) and 100 mol\% As40Se60, was made using high purity As (99.99999%, Furukawa Electric Ltd.), Se and S (both 99.999%, Matieron). As, and Se, were further purified by heating under vacuum (\(10^{-3}\) Pa) at 310 °C, and 270 °C, prior to batching (15–20 g samples), which was carried out inside a glove-box (0.1 ppm O2, ≤0.1 ppm H2O; MBraun). The silica glass ampoules (<1.0 ppm OH, ID/OD = 10 mm/14 mm, MultiLab.) in which the glasses were batched and melted, had been air-baked (1000 °C for 6 hours) and vacuum baked (1000 °C for 6 hours at \(10^{-3}\) Pa) prior to use. The ampoule was sealed under vacuum (~\(10^{-3}\) Pa), placed in a rocking furnace held statically, raised at 30 °C h⁻¹ to 250 °C and held isothermally for 2 h to ensure the chalcogens had fully melted. The furnace was then programmed to rock (±30° to horizontal axis) as the ampoule was heated to ~650 °C (30 °C h⁻¹) before being held isothermally with rocking for 12 h to achieve homogenisation. The furnace was then placed statically vertical and cooled at 60 °C h⁻¹ to ~570 °C. A final isothermal dwell of ~2 hours, to fine and refine the glass, was given before quenching the glass. The ampoules were quenched using a jet of nitrogen gas and subsequently annealed at \(T_g\) for 30 minutes, before cooling slowly to ambient, producing small glass rod samples.

Sample compositions were characterised using a PANalytical MiniPal4 to collect X-ray Fluorescence (XRF) spectra and Scanning Electron Microscopy, in back scattered mode, was used to collect Energy Dispersive X-ray (EDX) spectra. Density measurements were carried out using a Quantachrome Micropymometer, with helium as the displacement gas. Measurements of the glass transition temperature were carried out using a Perkin Elmer DSC7 Differential Scanning Calorimeter with water circulating chillers and an argon gas flow at the rate of 20 ml min⁻¹. For each measurement an initial “conditioning” temperature run was carried out by heating, and then cooling, the glass through the glass transition temperature, using a ramp rate of 10 °C min⁻¹ to ensure all glasses had the same thermal history. The sample was then heated again at 10 °C min⁻¹ to determine \(T_g\). The value for \(T_g\) was calculated using two different methods. The first method was to determine the point of maximum gradient, \(T_{g,max}\), taken as the maximum point in the derivative of the data. The second method was to calculate an “onset temperature”, \(T_{g,onset}\), using the intersection of two tangents fitted to the baseline, and the point of...
maximum gradient in the \( T_\text{g} \) feature. The first method gives a larger value for \( T_\text{g} \) than the second.

3.2. Z-scan measurements of \( n_2 \)

A standard Z-scan setup\(^{21}\) was used to measure the nonlinear optical absorption, \( \beta \), and nonlinear refractive index, \( n_2 \), of the seven As–S–Se glasses. The laser system was a Spectra Physics optical parametric generator-optical parametric amplifier combination with a Ti:sapphire laser and regenerative amplifier. The laser pulse width was \( \sim 120 \) fs and the corresponding spectral width was 12 nm. The measurements were performed using a laser of wavelength 1550 nm, pulse repetition rate of 1 kHz and a beam waist of 35 \( \mu \)m. Open and closed (20\% transmission) aperture scans were recorded simultaneously to measure the nonlinear coefficients. The samples were prepared by cutting a 10 mm diameter discs from the as-annealed glass rods. Both surfaces of the discs were polished to a 1 \( \mu \)m finish. To ensure the surfaces were parallel, the discs were then hot pressed under vacuum between tungsten carbide plates at a temperature \( \sim 40 \) °C above \( T_\text{g} \). For the measurement, the samples were mounted on a motorised translation stage and the movement and data collection were controlled by computer. Measurements were made at 2 energies (between 50 and 200 \( \mu \)J, chosen as appropriate for each glass composition) to confirm the reliability of \( n_2 \) calculated. Further details on the analysis of these samples are given by McCarthy.\(^{22}\)

3.3. Raman scattering

A Horiba LabRAM HR Raman system with a 660 nm Laser-Quantum Ignis laser (187 mW set to 25\% power) was used to collect Raman scattering data from the sample material. The samples were prepared by cutting a disc from the as-annealed glass rod and polished to a 1 \( \mu \)m finish. The data were collected using a 4 second exposure and accumulated over 10 scans. Repeat measurements at different points on each glass disc were made to ensure that the Raman spectra measured were representative of the whole glass. A background subtraction was performed on each spectrum before fitting was carried out.

3.4. \(^{77}\)Se NMR

\(^{77}\)Se NMR spectra were recorded at room temperature on an Avance III 600 Bruker spectrometer at a Larmor frequency of 114.497 MHz using a 2.5 mm magic angle spinning (MAS) probe, spinning at 15 kHz. A spin echo pulse sequence was used to acquire the broad line that results from the distribution of isotropic chemical shifts expected for the glassy state. The echo delay was synchronised with the MAS period and set to 66.67 \( \mu \)s. The \(^{77}\)Se \( \pi/2 \) pulse duration was 4 \( \mu \)s and the spectral width was 500 kHz. Acquisition commenced immediately after the \( \pi \) pulse in order to sample the whole echo to increase the signal-to-noise ratio. Given the low natural abundance of \(^{77}\)Se (7.58\%), up to 24 000 scans were required per spectrum, with a relaxation delay of 10 s. Chemical shifts are quoted relative to neat \((\text{CH}_3)_2\text{Se}\).

3.5. Neutron and X-ray total scattering

For total scattering using neutrons, the as-annealed glass rods were placed inside cylindrical vanadium containers, of inner diameter 10 mm, for the neutron scattering measurements. These vanadium containers were thin-walled, with a thickness of 25 \( \mu \)m, in order to reduce the magnitude of the experimental corrections required. The neutron diffraction patterns were measured using the GeM diffractometer\(^{23}\) at the ISIS pulsed neutron and muon source (Rutherford Appleton Laboratory, Oxfordshire, UK). The data were normalised using a calibration measurement of an 8 mm vanadium rod and were reduced and corrected for attenuation and multiple scattering using standard Gudrun\(^{24}\) and ATLAS\(^{25}\) software to give the distinct scattering spectra, \( I(Q) \). \( I(Q) \) was fitted with a quadratic, of the form \( A + BQ^2 \), at low \( Q \) to extrapolate the data to \( Q = 0 \).

The X-ray data were collected at the Diamond Light Source (Rutherford Appleton Laboratory, Oxfordshire, UK) using the I15 beamline optimised for the collection of PDF data. The photon energy used for the experiment was 73 keV. The samples were powdered and placed in 1 mm diameter silica capillaries, and the scattering was detected using a Perkin Elmer 2D flat panel 1621 detector. The 2D data were reduced to 1D using Fit2D\(^{26,27}\) and were corrected for detector attenuation and multiple scattering using standard Gudrun\(^{24}\) and ATLAS\(^{25}\). The 1D data were normalised using the Krogh-Moe and Norman method\(^{17,18}\) and corrected for attenuation and multiple scattering using GudrunX\(^{29}\) to produce the sharpened distinct scattering spectrum, \( I(Q) \). In both experiments, measurements of the empty instrument and empty container were made for data corrections. Once fully corrected \( I(Q) \) and \( I(Q) \) were Fourier transformed, using the Lorch modification function and a maximum momentum transfer, \( Q_{\text{max}} \), of 35 Å\(^{-1}\) for neutron and 25 Å\(^{-1}\) for X-ray data, to give the total correlation functions, \( T(r) \).

Predictions of the interatomic distances between nearest neighbour atoms were made using standard bond valence parameters.\(^{30}\) The calculated bond lengths (presented in Table 2) were used as starting parameters when fitting \( T(r) \). Simultaneous fitting to \( T^n(r) \) and \( T^Q(r) \) was carried out using NNXFit.\(^{31}\)

The neutron diffraction data, in both reciprocal- and real-space, are available from the ISIS Disordered Materials Database.\(^{35}\)

4. Results

4.1. Sample characterisation

Densities ranged from 3.19 ± 0.04 g cm\(^{-3}\) for As\(_{40}\)S\(_{60}\) to 4.55 ± 0.04 g cm\(^{-3}\) for As\(_{40}\)Se\(_{60}\) (Table 3). The linear variation in density

| Bond type | Coordination number 2.00 | Coordination number 3.00 | Coordination number 4.00 |
|-----------|-------------------------|-------------------------|-------------------------|
| As–S      | 2.11                    | 2.26                    | 2.37                    |
| As–Se     | 2.24                    | 2.39                    | 2.50                    |
| As–As     | 2.26                    | 2.41                    | 2.52                    |
| S–S       | 2.07                    | 2.22                    | 2.33                    |
| Se–S      | 2.25                    | 2.40                    | 2.51                    |
| Se–Se     | 2.36                    | 2.51                    | 2.62                    |
Table 3 Compositional data for the seven As–S–Se glasses, as determined by XRF and EDX analysis, with density, non-linear refractive index, $n_2$, non-linear absorption coefficient, $\beta$, and Se–S coordination number measurements, $n_{SeS}$.

| Nominal composition at% | XRF composition at% (±3) | EDX composition at% (±1) | Density g cm$^{-3}$ (±0.04) | $T_g$ Onset (±2) | Maximum gradient (±2) | $n_2$ ($\times 10^{-18}$ m$^2$ W$^{-1}$) (±0.5) | $\beta$ ($\times 10^{-12}$ mW$^{-1}$) (±0.2) | $n_{SeS}$ (±0.4) |
|------------------------|--------------------------|---------------------------|-----------------------------|-----------------|-----------------------|---------------------------------|---------------------------------|------------|
| As$_{40}$S$_{60}$       | As$_{40}$S$_{60}$         | As$_{40}$S$_{60}$         | 3.19                        | 198             | 209                   | 2.3                             | 0.53                             |            |
| As$_{40}$S$_{60}$Se$_{15}$ | As$_{40}$S$_{60}$Se$_{15}$ | As$_{40}$S$_{60}$Se$_{15}$ | 3.39                        | 190             | 200                   | 2.3                             | 1.0                              | 4.4        |
| As$_{40}$S$_{65}$Se$_{20}$ | As$_{40}$S$_{65}$Se$_{20}$ | As$_{40}$S$_{65}$Se$_{20}$ | 3.69                        | 192             | 204                   | 3.4                             | 2.7                              | 4.3        |
| As$_{40}$S$_{70}$Se$_{15}$ | As$_{40}$S$_{70}$Se$_{15}$ | As$_{40}$S$_{70}$Se$_{15}$ | 3.91                        | 189             | 199                   | 4.5                             | 2.0                              | 3.6        |
| As$_{40}$S$_{75}$Se$_{30}$ | As$_{40}$S$_{75}$Se$_{30}$ | As$_{40}$S$_{75}$Se$_{30}$ | 4.10                        | 188             | 199                   | 6.0                             | 3.7                              | 2.6        |
| As$_{40}$S$_{80}$Se$_{45}$ | As$_{40}$S$_{80}$Se$_{45}$ | As$_{40}$S$_{80}$Se$_{45}$ | 4.55                        | 170             | 176                   | 10.5                            | 2.8                              | 2.8        |

Fig. 1 (a) The density measurements for six of the samples in this study compared to literature values published by Cardinal et al. The density for As$_{40}$S$_{60}$ is not shown as the sample size was too small to allow an accurate measurement. (b) The glass transition temperatures for all seven glasses are compared to the literature values published by Vlcek et al. and Snapatin et al. in addition to those published by Višníc et al. and Wang et al. The different results are grouped depending on the definition of $T_g$ used.

with $x$, shown in Fig. 1a, suggests that there are no non-linear long range changes in the arrangement of the basic structural [As$_3$Ch$_4$] units as Se is substituted for S. $T_g$ measurements are shown in Fig. 1b, calculated by two methods. The values of $T_{g\text{onset}}$ and $T_{g\text{onset}}$ were found to vary non-linearly with the addition of As$_{40}$S$_{60}$Se$_{15}$. There was a decrease in $T_{g\text{onset}}$ of 8 °C on changing the glass composition from As$_{40}$S$_{60}$ to As$_{40}$S$_{65}$Se$_{15}$. However, as further selenium is added to the glass network, it was observed that $T_{g\text{onset}}$ value remain constant, within error, until a composition of As$_{40}$S$_{15}$Se$_{45}$ is reached. On changing the composition of the glass from As$_{40}$S$_{15}$Se$_{45}$ to one with no sulfur at all, As$_{40}$Se$_{60}$, a further drop in $T_{g\text{onset}}$ of 18 °C was observed.

The XRF and EDX results for glass composition were calculated using standardless calibration programs and are also given in Table 3. EDX measurements for each element were correct, to within 2 at%, at all times. This indicated that, within error, all the glasses had a stoichiometric composition As$_{40}$Ch$_{60}$. However, it should be noted that although the compositional analysis carried out by XRF is also consistent with the batched compositions, within error, the binary glass, As$_{40}$S$_{60}$, deviated most from the nominal compositions (3 at% excess As). It was expected that XRF would determine the composition of this glass with reasonable accuracy as it is comprise only two components. Therefore, while the batch compositions are assumed to be correct for subsequent analysis, the possibility that As$_{40}$S$_{60}$ was arsenic rich is considered in the discussion.

4.2. Raman

The Raman spectral bands, shown in Fig. 2a, are centred at 345 cm$^{-1}$ for As$_{40}$S$_{60}$, characteristic of the vibration of As–S bonds, and 230 cm$^{-1}$ for As$_{40}$Se$_{60}$, characteristic of As–Se bond vibrations. The mixed chalcogen samples were composed of As–S and As–Se bonds and so exhibit both peaks. The effect of oscillator strength is shown in Table 4. The area of the As–S peak measured for As$_{40}$S$_{60}$ is ~50% of that measured for the As–Se peak in As$_{40}$Se$_{60}$. As a result, the relative intensities of the two peaks cannot be used directly to determine the relative amounts of As–S and As–Se bonds in the mixed chalcogen glasses. Fig. 2b shows the percentage variation in intensity of the As–S and As–Se peaks with $x$. The intensity of the As–Se peak, as Se was replaced by S, varied linearly, but the As–S peak intensity did not. A linear trend line fitted to the As–S peak intensities measured for $x = 0$–67 predicted that a Raman spectrum for a glass consisting of 85 mol% As$_{40}$Se$_{60}$ would exhibit no As–S peak.

4.3. $^{77}$Se NMR

To check the consistency of the selenium environment with composition, $^{77}$Se NMR spectra were collected for As$_{40}$S$_{60}$ and As$_{40}$S$_{40}$Se$_{20}$ ($x = 100$ and 33) glasses. Gaussian peak fits to the spectra, shown in Fig. 3, were fixed at 380 ppm and the NMR spectra for both glasses could be simulated well with a single peak. A peak with a chemical shift of 380 ppm arises from Se atoms bonded to two arsenic atoms.
4.4. Z-scan measurements of $n_2$

Z-scan measurements of the non-linear refractive index, $n_2$, and non-linear absorption, $\beta$, are given in Table 3. There is a slow change in $n_2$ in sulfur rich glasses, varying from $2.3 \pm 0.5 \times 10^{-18}$ $\text{m}^2 \text{W}^{-1}$ to $4.5 \pm 0.5 \times 10^{-18}$ $\text{m}^2 \text{W}^{-1}$ as the composition is changed from $\text{As}_{40}\text{S}_{60}$ to $\text{As}_{40}\text{S}_{30}\text{Se}_{30}$ ($x = 0$–50). However, as the selenium content was increased, the rate of change accelerated, reaching a value of $10.5 \pm 0.5 \times 10^{-18}$ $\text{m}^2 \text{W}^{-1}$ for $\text{As}_{40}\text{Se}_{60}$ ($x = 100$). Literature measurements for $n_2$ (ref. 3 and 4) are quoted relative to silica making direct comparison of values of $n_2$ difficult. To facilitate a comparison, Fig. 4a shows the values for $n_2$ measured by this study, along with those measured by Cardinal et al.,$^3$ Harbold et al.,$^4$ and Wang et al.,$^{32}$ relative to the values of $n_2$ measured for $\text{As}_{40}\text{Se}_{60}$. The relative change with $x$ observed by all three studies is similar, with $n_2$ for $\text{As}_{40}\text{Se}_{60}$ being 4.0–4.5 times larger than $\text{As}_{40}\text{S}_{60}$. Fig. 4b shows that the variation in $n_2$ with $x$ is not linear, but can be fitted with a second order polynomial equation.

Table 4 The areas of As–Se and As–S bands in the Raman spectra obtained for each composition using a multi-peak fit to the peak manifold. Results are given as a % of the peaks observed in the end member glasses.

| Nominal composition at% | As–Se peak area as % of $\text{As}_{40}\text{Se}_{60}$ spectrum | As–S peak area as % of $\text{As}_{40}\text{S}_{60}$ spectrum |
|-------------------------|-----------------------------------------------------------|-----------------------------------------------------------|
| $\text{As}_{40}\text{S}_{60}$ | 0.0 | 100.0 |
| $\text{As}_{40}\text{S}_{45}\text{Se}_{15}$ | 22.3 | 65.2 |
| $\text{As}_{40}\text{S}_{30}\text{Se}_{20}$ | 31.1 | 59.0 |
| $\text{As}_{40}\text{S}_{20}\text{Se}_{30}$ | 46.0 | 40.0 |
| $\text{As}_{40}\text{S}_{15}\text{Se}_{45}$ | 61.9 | 20.1 |
| $\text{As}_{40}\text{S}_{30}\text{Se}_{15}$ | 70.8 | 13.3 |
| $\text{As}_{40}\text{S}_{60}$ | 100.0 | 0.0 |

Fig. 2 (a) The Raman spectra for seven bulk chalcogenide glasses, collected at 660 nm. Arrows indicate the position of the shoulders on the As–S peak. (b) The relative areas of the As–S (blue) and As–Se (red) peaks with composition. The solid lines show the best fit through all the data points and the blue dashed line is a linear fit to the As–S peak intensities for $x < 67$. (c) A more detailed comparison of the changes in the As–Se peak shape with composition (key matches that in a). Each spectrum is scaled by the amount of $\text{As}_{40}\text{Se}_{60}$ present in the sample, $x$.

Fig. 3 $^{77}\text{Se}$ NMR spectra acquired for the $\text{As}_{40}\text{Se}_{60}$ and $\text{As}_{40}\text{S}_{60}$ samples. The solid lines are the experimental data, and the dashed lines are the peak fits.

Fig. 4 (a) The variation in non-linear optical response of the series of $x\text{As}_{40}\text{Se}_{60}(100 - x)\text{As}_{40}\text{S}_{60}$ glasses as measured in this study at 1.55 $\mu\text{m}$ wavelength (black squares). The results are compared to literature results (right y axis) from Cardinal et al.$^3$ at 1.6 $\mu\text{m}$ (red squares), Harbold et al.$^4$ at 1.55 $\mu\text{m}$ and Wang et al.$^{32}$ at 1.55 $\mu\text{m}$. To allow direct comparison, the results are normalised to the $n_2$ value for $\text{As}_{40}\text{S}_{60}$ ($x = 100$). Literature measurements for $n_2$ (ref. 3 and 4) are quoted relative to silica making direct comparison of values of $n_2$ difficult. To facilitate a comparison, Fig. 4a shows the values for $n_2$ measured by this study, along with those measured by Cardinal et al.$^3$, Harbold et al.$^4$ and Wang et al.$^{32}$ relative to the values of $n_2$ measured for $\text{As}_{40}\text{Se}_{60}$. The relative change with $x$ observed by all three studies is similar, with $n_2$ for $\text{As}_{40}\text{Se}_{60}$ being 4.0–4.5 times larger than $\text{As}_{40}\text{S}_{60}$. Fig. 4b shows that the variation in $n_2$ with $x$ is not linear, but can be fitted with a second order polynomial equation.
4.5. Neutron and X-ray scattering

The distinct scattering patterns, measured using neutrons, \( I^N(Q) \), and X-rays, \( I^X(Q) \), are shown in Fig. 5 (except As\(_{40}\)Se\(_{60}\), which was not measured). All diffraction data exhibited a shift in the position of peaks to shorter \( Q \) with increasing selenium content. The exception was the position of the pre-peak at 1.26 Å\(^{-1} \), which was unchanged (±0.3 Å\(^{-1} \)). The shift in peak position to shorter \( Q \) indicated an expansion of the glass network to accommodate the increase in average chalcogen size. The corresponding correlation functions, \( T^N(r) \) and \( T^X(r) \), are shown in Fig. 6. There were two well defined peaks at \( \sim 2.35 \) Å and \( 3.55 \) Å in the correlation functions; the former arose from As–Ch bonds and the latter from the shortest, non-bonded, distances for As–As/Ch–Ch. Some of these short As–As and Ch–Ch distances arose from the next nearest neighbours within an [As\(_n\)Ch\(_3\)] unit (Ch–As–Ch) or between two units (As–Ch–As), but an examination of the structure of crystalline As\(_{40}\)Se\(_{60}\) (ref. 33) indicated that up to 50% of the contribution to \( n_{AsAs} \) may not be caused by As–Ch–As through-bonding, but rather from the close packing of structural units that are not directly connected.

Fig. 7 shows simultaneous fits to the first peak in \( T^N(r) \) and \( T^X(r) \) for the \( xAs_{40}Se_{60}(100 - x)As_{40}Se_{60} \) glass system. The fits required two peaks in order to model the As–S and As–Se bonds. The positions, widths and areas of the peaks were allowed to vary freely and the fit parameters are given in Table 5. The As–Se and As–S peak positions were all within 0.01 Å of the bond lengths for As–S and As–Se (2.27 and 3.62 Å) and As–S and As–Se (2.41 and 3.42 Å) were consistent with Ch–As–Ch bond angles of \( \sim 100^\circ \), which is in good agreement with that expected for the trigonal pyramid environment of a three-coordinated lone-pair ion.34

5. Discussion

5.1. Density, glass transition temperature and composition

The measured densities for glasses in the system \( xAs_{40}Se_{60}(100 - x)As_{40}Se_{60} \) shown in Fig. 1a agreed, within error...
A change in definition of Snopatin et al. suggests that there may have been a small excess of arsenic, suggesting that there may have been a small excess of arsenic.

Table 5 The parameters obtained from simultaneous fits to the As–Ch peaks in X-ray and neutron T0/cs. The bond lengths, r, peak widths, σ, and coordination numbers, nAsCh, are given, along with the expected coordination numbers, for each sample. The two peak fit for As40S60 is treated as both two As–S distances, and As–S and an As–As distance.

| Composition   | Correlation       | r (Å)    | σ (Å)   | nAsCh   | Calculated nAsCh |
|---------------|-------------------|----------|---------|---------|------------------|
| As40S60       | As–S              | 2.271(1) | 0.0585(2) | 2.86(3) | 3                |
|               | As–S (As–As)      | 2.467(3) | 0.029(1)  | 0.11[3] (0.09[3]) |               |
|               | As–S (As–X) total | 2.97[5] (2.95[5]) |         |         |                   |
| As40S40Se15   | As–S              | 2.273(2) | 0.061[1]  | 2.24(3) | 2.25             |
|               | As–Se             | 2.424(2) | 0.074[1]  | 0.74(3) | 0.75             |
|               | As–Ch total       | 2.98(3) |          |         |                   |
| As40S40Se20   | As–S              | 2.277(2) | 0.066[1]  | 1.97(3) | 2                |
|               | As–Se             | 2.419(2) | 0.069[1]  | 0.97(3) | 1                |
|               | As–Ch total       | 2.94(3) |          |         |                   |
| As40S40Se30   | As–S              | 2.279(2) | 0.061[1]  | 1.49(3) | 1.5              |
|               | As–Se             | 2.416(2) | 0.065[1]  | 1.48(3) | 1.5              |
|               | As–Ch total       | 2.97(3) |          |         |                   |
| As40S40Se40   | As–S              | 2.277(2) | 0.068[1]  | 0.95(3) | 1                |
|               | As–Se             | 2.417(2) | 0.069[1]  | 1.95(3) | 2                |
|               | As–Ch total       | 2.90(5) |          |         |                   |
| As40S60       | As–Se             | 2.412(1) | 0.0659(2) | 2.93(3) |                   |

(±0.04 g cm⁻³), with previously work by Cardinal et al.³ and Snopatin et al.⁵ and are a first indication that the glass compositions did not deviate from those batched. In contrast, the measurements for the glass transition temperature do not agree with those reported by Cardinal et al. and Snopatin et al. for the same samples.

When comparing the glass transition temperatures, Tg, measured in this study with those in the literature (Fig. 1b), it is important to consider how Tg is defined by different researchers. Cardinal et al. clearly state that they define Tg as the point in the region of glass transition with the maximum gradient.³ Data published by Vlček et al.⁶ yield very similar results to Cardinal et al. and, though it is not explicitly stated, we have assumed that they used a similar definition of Tg. In contrast, Tg values measured by Abdel-Moneim were reported to be calculated using the onsets Tg definition.¹³ The results published by Snopatin et al. for comparable compositions are consistently lower in temperature than all other data sets,⁵ but the equivalence between the g values reported for As 40Se60 in this study was closer to those batched. In contrast, Tg values for these glasses were closer to that of As40S60 than As40Se60 indicating that the glass transition temperature for glasses in a compositional range from As40S40Se40 to As40Se60 have a considerable impact on the interpretation of the data. An analysis of the glass transition temperatures obtained in this study suggest that Tg does not vary linearly in mixed anion glasses, but is instead relatively constant across a wide compositional range.

The invariance in the glass transition temperature for glasses in a compositional range from As40S40Se15 to As40Se60 suggests that the mixed phase glasses maintain a similar structural motif. Furthermore, the Tg results for these glasses were closer to that of As40S60 than As40Se60 indicating that the glass relaxation behaviour is more similar to the sulfide glass than the selenide. Careful analysis of the composition and homogeneity of the glasses is vital to ensure that this interpretation is accurate. The presence of phase separation in the As40Se60 glass, for example, would result in a mix of As–As and Se–Se bonds that could affect Tg. Literature indicates that moving away from the stoichiometric composition, and therefore increasing the number of homopolar bonds, lowers Tg.³⁶,³⁷ However, it should be noted that a review of Tg temperatures for As40Se60, reported in the Sciglass software, gave a wide variation in results, ranging from 150 °C to 200 °C, and yielded a mean Tg of 181 °C with a standard deviation of 11 °C. The “anomalous” Tg values reported for As40Se60 in this study were both within this error range.

Compositional analyses of all of the glasses agreed with the batched compositions, within the error of the techniques used (Table 3). The composition for As40S60 determined using XRF (Table 3) was most disparate from the batched composition, suggesting that there may have been a small excess of arsenic in the glass. If this were true, the glass would no longer be
comprised of only [AsS₃] units, but must have included a number of As–As bonds. The As₄₀Se₆₀ glass, which has an anomalous glass transition temperature, is reported to have a composition of As₄₀Se₆₀, and so is considered to be stoichiometric in all analysis.

5.2. Short range order and homopolar bonds

The glass compositions used in this study have been designed so that, in an ideal glass, each As atom was bonded to three Ch atoms, and each Ch is bonded to two As atoms. Therefore, if homopolar bonds were present, they would have arisen from either an excess of As (as suggested by XRF compositional analysis for As₄₀Se₆₀), or from phase separation resulting in As rich and Ch rich regions (see, for example, work by Petri et al.²⁹ and Georgiev et al.⁴⁰). Phase separation can be intrinsic to the glass, or arise from a lack of inhomogeneity in the glass melt. The structural information gained from Raman spectroscopy, ⁷⁷Se NMR and neutron and X-ray total scattering were therefore interrogated for evidence of homopolar bonds before stoichiometry (As₄₀Ch₆₀) was accepted.

5.2.1. Evidence for homopolar bonds in As₄₀S₆₀. XRF analysis of the As₄₀S₆₀ glass indicated that the true composition of the glass was As₄₃S₇. However, this result was not considered reliable because the Raman spectrum in Fig. 2a and T⁹(r) in Fig. 6a did not match the published results for glasses with compositions of ~As₁₂S₅₈.⁹,¹³,¹⁴ Fig. 2a shows that a peak manifold, centred at ~350 cm⁻¹, was the dominant contribution to the Raman spectrum of As₄₀S₆₀, as expected for As–S bonds. However, there were also weak bands observable at 188, 233 and 498 cm⁻¹ that may indicate the presence of a very small number of As–As and S–S bonds in the glass.²,¹³,¹⁴ For example, peaks at ~500 cm⁻¹ in the Raman spectra have been assigned by Golovchak et al. to sulfur dimers of the form As–S–S–As.⁹ There was no evidence in the neutron or X-ray correlation functions (Fig. 6a) of S–S peaks, which would have contributed a peak in T(r) at 2.04 Å (see Table 2 and ref. 13), but it should be noted that sulfur scatters both X-ray and neutrons very weakly in contrast to arsenic.

Georgiev et al.⁴⁰ attributed the weak features, observed in the Raman spectrum of As₄₀S₆₀, at 188 and 233 cm⁻¹ to the vibrations of [As₅S₃] molecules, which are comprised of As–S and As–As bonds. Using this assignment, and the intensity of the two bands, they estimated that 7% of the As bonds in bulk As₄₀S₆₀ glass were homopolar.⁴⁰ The correlation functions obtained using neutrons and X-rays, T³(r) and T⁶(r), were carefully analysed to determine where this concentration of As–As bonding is reasonable.

The simultaneous fits to the correlation functions, T³(r) and T⁶(r), for As₄₀S₆₀ (Fig. 7) required two peaks to model the distribution of bond lengths fully. The parameters for the fit are given in Table 5. The peak positions, 2.271 ± 0.001 Å and 2.467 ± 0.001 Å, could be interpreted as arising from a distribution of As–S bond lengths, or from As–S and As–As bonds, respectively (see Table 2 and ref. 13). A consideration of the total arsenic coordination number does not aid in determining which interpretation is correct because both yield a total coordination number of 3, within the error of the experiment. The width of the peak fitted at 2.467 Å was 0.029 ± 0.001 Å. This is significantly narrower than the widths of the As–S or As–Se peaks given in Table 5 (~0.06 Å) and indicates that the peak is not a real measure of a correlation. Furthermore, the second peak is also at a distance that is more than 0.06 Å shorter than the As–As bond length measured using neutrons¹³ and K edge As EXAFS⁴² in As rich glasses. Based upon this interpretation, it can be proposed that the slight asymmetry of the As–S peak shape in the neutron correlation function (Fig. 6) arises from a distribution of As–S bond lengths in the glass. Similar asymmetric bond distributions have been observed in correlation functions for oxide glasses containing lone pair cations, such as As, when a Q max ≥ 30 Å⁻¹ is used for the Fourier transform (i.e. [PbO₃]¹³ and [SbO₃]¹⁴). For example, a similar two peak fit to the Pb–O distribution in PbO–Al₂O₃ glasses yielded Pb–O peak positions of ~2.25 and 2.47 Å. In contrast, the presence of small peaks in the Raman spectra at ~200 cm⁻¹ supports the interpretation that the shoulder arises from a small number of As–As distances in the glass. An As–As coordination number of 0.09 ± 0.03 indicates that between 2 and 4% of arsenic bonds are homopolar, slightly less than the number predicted by Georgiev et al.⁴⁰ In summary, the measured coordination number of arsenic is close to the expected value of three, and there is some evidence for a small number of As–As bonds. However, the presence of homopolar bonds at these concentrations would have a negligible effect on further analysis and the glass shall be considered as stoichiometric and comprised solely of [AsS₃] in the remainder of the discussion. For a more detailed study of As–As bonds in As₄₀S₆₀ careful fitting of T⁸(r) for a series of AsₓS₁₀₀–ₓ glasses is required to establish robust parameters for S–S, As–S and As–As peaks.

5.2.2. Evidence for homopolar bonds in Se containing glasses. The close similarity of atomic mass and radius for As and Se makes the determination of homopolar Se–Se bonds in the presence of As–Se bonds more difficult than identifying S–S bonds in the presence of As–S. Bond valence parameters⁴⁰ were used to calculate the expected bond lengths for an arsenic atom bonded to three Se or As atoms, and a selenium atom bonded to two Se atoms. The results are given in Table 2 and the three distances, differing by less that 0.05 Å, would not resolvable in a neutron or X-ray correlation function. Furthermore, the vibrational frequencies of As–Se, Se–Se and As–As bonds are also unresolvable by Raman spectroscopy. Therefore, an alternative determination of the local environment of Se in the arsenic selenide glasses was used. ⁷⁷Se NMR studies of arsenic selenide based glasses carried out by Bureau et al.⁷,³ and Delaizir et al.⁴⁵ have identified three chemical shifts, 380, 560 and 820 ppm, in ⁷⁷Se NMR spectra. These shifts have been assigned to As–Se–As, As–Se–Se and Se–Se–Se environments, respectively. The ⁷⁷Se NMR spectra for As₄₀Se₆₀ and As₄₀Se₆₀Se₂₀, shown in Fig. 3, have been fitted with a single peak centred at 380 ppm, indicating that both glasses are primarily comprised of As–Se–As environments. The glasses in this study are therefore concluded to have a high degree of chemical order; Se is predominantly found in [AsCh₃] units, with few or no Se–Se bonds in the amorphous As₄₀Se₆₀ and As₄₀Se₆₀Se₂₀. This conclusion is
supported by Bureau and colleagues who used a $^{77}$Se NMR 2D Carr–Purcell–Meiboom–Gill (CPMG) pulse program to measure Se environments in As–S–Se glasses.46

5.3. Changes in network connectivity with composition

5.3.1. Raman. The intensities of the As–S band (270–420 cm$^{-1}$) in the Raman spectra (Fig. 2a) for the seven glass compositions are given in Table 4. The fractional change in intensity of the As–S band with $x$ is shown in Fig. 2b with a linear fit to the results. The fit indicated that there should be no contribution to the Raman spectra from an As–S peak when $x \geq 85$ in $x$As$_{40}$Se$_{60}$ (100 $-$ $x$)As$_{40}$S$_{60}$ glasses. However the intensity of the As–S band at $\sim 350$ cm$^{-1}$ was distorted by the presence of shoulders to either side of the main peak. These are labelled with black arrows in Fig. 2a. Wagner et al.11 attributed these shoulders to vibrations between $[\text{AsS}_3]$ units (i.e. $\text{S}_x\text{As}_y\text{S}_z\text{As}_w$). The As–S band in the experimental Raman spectra, shown in Fig. 2a, narrowed as S was replaced with Se until a composition of $x$As$_{40}$S$_{60}$ was reached, indicating a decrease in intensity of the shoulder bands; the widths of the As–S peaks in Raman spectra for $x$As$_{40}$S$_{20}$Se$_{40}$ and $x$As$_{40}$S$_{15}$Se$_{45}$ ($x = 67$ and 75 respectively) showed no change (Fig. 2a). By using two lines to fit the As–S peak intensities in Fig. 2b, one to the $x = 0$–50 peak intensities and one to the $x = 67$–100 values (where $x = 100$ is an As$_{40}$Se$_{60}$ glass with no As–S peak), a more accurate measure of the sulfur concentration at which $[\text{AsS}_3]$ units are no longer directly connected could be obtained from the intersection of the two lines. This occurred at $x = 70$ (As$_{40}$S$_{18}$Se$_{42}$). If the interpretation that the shoulders in the spectral band arise from bonded $[\text{AsS}_3]$ units is correct, they would not be expected for compositions beyond $\sim 33\%$ As$_{40}$Se$_{60}$. At this composition, the average arsenic unit would be $[\text{AsS}_2\text{Se}]$. The persistence of intensity in the shoulders to $\sim 70\%$ As$_{40}$Se$_{60}$ suggests that either the interpretation is incorrect, or that $[\text{AsS}_3]$ units are present in the glass in much higher concentrations than would be expected by random chalcogen mixing.

Fig. 2c shows the As–Se peak manifold in detail. The Raman spectra have been scaled by the atomic fraction of Se present in the glass to allow a direct comparison the peak intensities and shape. If all Se atoms are bonded to As, the area of the bands in Fig. 2c should be invariant with composition. The linear relationship between $x$ and peak area in Fig. 2b shows this to be true. The average $[\text{AsCh}]_x$ structural unit in an As$_{40}$Se$_{60}$ glass changes with the addition of As$_{40}$Se$_{60}$ from $[\text{AsS}_3]_x$ via the formation of $[\text{AsSe}_3]_x$ and $[\text{AsS}_3]_x$ units as Se replaces S in the glass network, until all As atoms are in $[\text{AsSe}_3]_x$ units. A clear shift in the position of the maximum intensity of the As–Se band to lower wave numbers was observed as Se was added to the glass. Freitas et al.11 and Li et al.10 have previously attributed this red shift to a change in the next nearest neighbours for the selenium atoms. It should be noted that the peak manifolds at $\sim 250$ cm$^{-1}$ for As$_{40}$Se$_{60}$ and As$_{40}$S$_{60}$Se$_{20}$ ($x = 25$ and 33) were similar, indicating that there was little change in the local Se environment at these two compositions.46

5.3.2. Ch–Ch coordination numbers determined from correlation functions. It is possible to estimate the correlation function for a ternary As–S–Se glass using a weighted sum of the correlation functions for the binary glasses. A weighted function can be calculated by summing $T^N(r)$ for the two binary glasses, $x$As$_{40}$Se$_{60}$ and $x$As$_{40}$S$_{60}$ in proportions to match the composition of each ternary glass, $x$As$_{40}$Se$_{60}$($100 - x$)As$_{40}$S$_{60}$. However, this sum is only an approximation for $T^N(r)$ for a ternary glass because $x$As$_{40}$Se$_{60}$ and $x$As$_{40}$S$_{60}$ are solely composed of $[\text{AsS}_3]_x$ and $[\text{AsSe}_3]_x$ units. Therefore, even if a weighted sum of the $\text{As}\cdot\text{As}$ and $\text{As}\cdot\text{Ch}$ correlations in $x$As$_{40}$Se$_{60}$ and $x$As$_{40}$S$_{60}$ glasses were a good approximation for the corresponding correlations in the mixed glasses, it cannot properly account for the Ch–Ch distances that arise from the presence of $[\text{AsS}_2\text{Se}]$ and $[\text{AsS}_3\text{Se}]$ units. Differences between the experimental correlation function and the weighted sum would become apparent in the region of the second peak in $T^N(r)$, where the weighted Ch–Ch sum will over estimate the number of Se–Se and S–S correlations and neglect Se–S correlations. As the scattering length of S is significantly less than Se (Table 1), the weighted sum will have a more intense contribution from Ch–Ch correlations than is measured experimentally.

The experimental neutron correlation functions, $T^N(r)$, for four ternary glasses are shown in Fig. 8 with the calculated weighted correlation functions. There was excellent agreement between the experimental data and weighted sum over the distance ranges 0–3.2 Å and 4.25–5.4 Å. For each glass sample, the difference between the experimental correlation function and the calculated weighted sum in the region of the first peak in $T^N(r)$ was of the same order as the noise in the low $r$ region. The first significant peak in the residual, attributed to Se–S distances, was observed at $\sim 3.75$ Å and is narrow and well defined. The $x$As$_{40}$S$_{60}$Se$_{20}$ glass exhibits the largest residual, because this is the composition where the potential for mixing...
is maximised. Fig. 9 compares the differences between the experimental Se ··· S correlation and a calculated partial correlation for Se ··· Se distances (calculated from the reported crystal structures using the XTAL program47) in crystalline As₄₀Se₆₀. There is close agreement between the first peak maximum for this simulation (3.71 Å) with the residuals (3.75 Å) and a general agreement in the shape of the residuals with the Se ··· Se partial correlation function over the range of 0–10 Å. This supports the hypothesis that the differences shown in Fig. 9 arose from a Ch ··· Ch correlation.

\( n_{\text{SeS}} \) values for each glass containing selenium were calculated by integrating the area of the peaks at \( \sim 3.75 \text{ Å} \) in Fig. 9 and using eqn (11). The results are given in Table 3 and shown in Fig. 4b. If the Se and S anions were evenly distributed throughout the glass, a linear variation would be expected. A linear fit was made to \( n_{\text{SeS}} \) for glasses containing \( \geq 50 \text{ mol\% As₄₀Se₆₀} \), and is shown as a red dashed line in Fig. 4b. An extrapolation of this linear fit to \( x = 0 \) indicated that a single Se atom in an arsenic sulphide glass would have 7.3 sulfur neighbours, \( n_{\text{SeS}} = 7.3 \). This is close to \( n_{\text{SeSe}} = 7.33 \), calculated from the crystal structure of crystalline As₄₀Se₆₀.33 when using a maximum Se ··· Se distance of 4.0 Å. The good agreement between the extrapolated value for \( n_{\text{SeS}} \) in the As₄₀S₆₀ and As₄₀Se₆₀ in crystalline As₄₀Se₆₀ indicated that the change in \( n_{\text{SeS}} \) was directly dependent on the total number of sulfur atoms in the glass. Therefore, selenium and sulfur atoms were randomly mixed when \( > 50\% \) of Ch atoms are selenium. However, Fig. 4b shows that values for \( n_{\text{SeS}} \) deviated from this linear correlation when \( < 50\% \) of the chalcogen atoms in the glass were selenium. To achieve a good fit to all values of \( n_{\text{SeS}} \) a second order polynomial was required. The intersection of this new fit with \( x = 0 \) occurred at 4.6. Such a markedly non-linear variation in \( n_{\text{SeS}} \) with composition suggested that either there was a preference for like atoms to cluster when there was a high concentration of sulfur in the glass, or that the average number of anions around the Se atom decreased from 7.3 to 4.6 as the amount of sulfur in the glass was increased from 0 to 60 at%. Conclusive evidence in favour of either suggestion is lacking. However, if the total number of Ch atoms around selenium decreased as sulfur was added to the glass, a corresponding expansion of the glass network would be expected. Instead, the shift in peak positions in \( I^2(\mathbf{Q}) \), shown in Fig. 5a, indicated that a network contraction occurs with increased sulfur content.

The interpretation of the Raman spectra in Section 4.3.1 supports the suggestion that \( n_{\text{SeS}} \) varies non-linearly with \( x \). The shape of the As–Se band in the Raman spectra is strongly influenced by the surrounding environment. However, the shape of the As–Se peaks in Fig. 2c for As₄₀S₃₀Se₁₀ and As₄₀Se₆₀Se₂₀ are very similar. This suggests that the local environment of the As–Se bonds was not changing as fast as would be expected if S and Se atoms were mixing randomly in the glass. Furthermore, the non-linear change in the intensity of the As–S peak in the Raman spectra suggested a preference for [AsS₃] units to persist in the glass to higher concentrations than random mixing would allow. In summary, the combination
of Raman and total scattering data indicates that there may be preferential clustering of selenium and sulfur in $x\text{As}_{40}\text{Se}_{60} - (100 - x)\text{As}_{40}\text{S}_{60}$ glasses, when there is a high concentration of the sulfur in the glass. As a selenium atom has a radius that is $\sim 20\%$ larger than sulfur, this proposed clustering could be driven by an inability for a network of $[\text{AsS}_3]$ units to accommodate easily the larger Se atoms.

5.4. A comparison of glass structure with physical properties

The results presented here, coupled with the literature, show that $n_2$ varies non-linearly with composition in this glass system. Several theories have been put forward to explain the structural changes that underpin this non-linear response. Nasu et al. related an increase in non-linear optical susceptibility, $\chi^{(3)}$ (related to $n_2$) to a denser glass network, highlighting that the low $\chi^{(3)}$ measured for GeSe$_4$ corresponds to a high molar volume. However, this is not consistent with the glass system studied here, where the molar volume and $n_2$ both increased as selenium atoms were doped into the network. More recent work by Harbold et al. showed that there is a relationship between the band gap energy, $E_g$, and $n_2$. $n_2$ was shown to increase rapidly when the photon energy of the laser was $> 0.5E_g$. Sanghera et al. explicitly related $n_2$ with this normalised photon energy. However, it should be noted that the band gap energy and the onset of the absorption edge, $\lambda_{\text{gap}}$, also showed evidence of non-linear variation with composition.

Therefore, we suggest that the same changes in the glass structure affect both $\lambda_{\text{gap}}$ and $n_2$. Petit et al. showed an inverse relationship between the density of lone-pairs in the glass and $n_2$. However, this is in direct conflict with the accepted understanding that lone-pairs enhance $n_2$ in heavy metal oxide glasses (see ref. 52 for example). Harbold et al. highlighted this inconsistency and demonstrated that it is possible to maintain constant $\lambda_{\text{gap}}$ and $n_2$ values for selected glasses in the Ge–As–Se system that exhibit different lone pair densities.

From the brief discussion above, it is clear that the relationship between $n_2$, $\lambda_{\text{gap}}$, and glass structure is complicated and remains unresolved. The results of this study suggest an alternative structural origin for the non-linear variation in $n_2$ with $x$ in the $x\text{As}_{40}\text{Se}_{60} - (100 - x)\text{As}_{40}\text{S}_{60}$ glass system. Fig. 10 shows a linear relationship between $n_2$ and $n_{\text{SeS}}$. We therefore propose that the variation in $n_2$ in this series of glasses is related to the arrangement of the chalcogen atoms in the glass. The variation in $n_{\text{SeS}}$ only varied linearly, indicating a well-mixed glass, when $x > 0.5$. For glasses containing low concentrations of Se, there was evidence that the Se and S atoms preferentially cluster, which may result in a sublinear change in the non-linear refractive index of the glass. Work by Maklad et al. supports this interpretation as follows. Careful observation of I absorptions in the $x\text{As}_{40}\text{Se}_{60} - (100 - x)\text{As}_{40}\text{S}_{60}$ glass system indicated that the expected vibrations from $[\text{AsS}_3\text{Se}]$ and $[\text{AsSSe}_2]$ units are missing from an $x = 40\%$ $\text{As}_{40}\text{Se}_{60}$ glass, and are only weakly present in an $x = 93\%$ $\text{As}_{40}\text{S}_{60}$ glass. The lack of vibrations from mixed units has been interpreted as a sign of non-random distribution of the chalcogens in these glasses and suggests that $[\text{AsS}_3]$ and $[\text{AsSSe}_2]$ units form in preference to mixed units such as $[\text{AsS}_2\text{Se}]$, when Se is present in low concentrations. This is in good agreement with the structural results presented here, which suggests that full mixing of the chalcogen atoms does not occur until $\sim 50\%$ $\text{As}_{40}\text{Se}_{60}$ is present. Furthermore, the thermal analysis suggests that the mixed chalcogen glasses retain a more sulfide-like relaxation behaviour across the compositional range.

6. Conclusions

A series of $x\text{As}_{40}\text{Se}_{60} - (1 - x)\text{As}_{40}\text{S}_{60}$ glasses, where $x = 0, 25, 33.3, 50, 66.7, 75$ and $100\%$ $\text{As}_{40}\text{Se}_{60}$, was investigated. Values of $n_2$, measured using Z-scan, were found to vary between $2.3(5) \times 10^{18}$ m$^2$ W$^{-1}$ for $\text{As}_{40}\text{S}_{60}$ to $10.5(5) \times 10^{18}$ m$^2$ W$^{-1}$ for $\text{As}_{40}\text{Se}_{60}$ and show a non-linear variation with composition that was modelled using a second order polynomial equation. These results are in reasonable qualitative agreement with the trends in $n_2$ reported elsewhere. Through careful analysis of total correlation functions, obtained using neutrons and X-rays, it has been shown that there is a lower than expected number of sulfur atom nearest neighbours around a selenium atom in glasses containing $\leq 50\%$ $\text{As}_{40}\text{Se}_{60}$. The variation of $n_{\text{SeS}}$ with $x$ requires a second order polynomial fit to the data and this non-linear variation is attributed to the clustering of chalcogen atoms in the glass. The observed variation in $n_{\text{SeS}}$ was shown to have a linear relationship with $n_2$, and it is suggested that the spatial distribution of chalcogen atoms in the glass is a structural factor contributing to the non-linear changes in $n_2$ with composition. To be able to investigate the medium range structure of this As–S–Se glasses and determine the partial correlation functions for Ch · · · Ch correlations more fully structural modelling methods, such as Molecular Dynamics and Reverse Monte Carlo, will be used.

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