Dynamical behavior of Borospherene: A Nanobubble

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The global minimum structure of borospherene (B_{40}) is a cage, comprising two hexagonal and four heptagonal rings. Born-Oppenheimer Molecular Dynamics simulations show that continuous conversions in between six and seven membered rings take place. The activation energy barrier for such a transformation is found to be 14.3 kcal·mol\(^{-1}\). The completely delocalized \(\sigma\)- and \(\pi\)-frameworks, as well as the conservation of the bonding pattern during rearrangement, facilitate the dynamical behavior of B_{40}. B_{40} is predicted to act as a support-free spherical two-dimensional liquid at moderate temperature. In other words, B_{40} could be called as a nanobubble.

Is it possible to build a buckyball comprising only boron atoms? Boron is an electron deficient atom with only three valence electrons. So, a perfect boron buckminsterfullerene (B_{60}) is not expected to be stable owing to the absence of the fourth electron, which is essential for the \(\pi\)-stabilization of a spherical shell. A fullerene-like boron cluster, B_{80}, was predicted in silico, which has structural similarity with C_{60} but with an additional boron atom at the center of each hexagon. However, several unsymmetrical B_{80} structures were found to be more stable than the fullerene-like geometry.

Quite recently, Zhai et al.7 reported the experimental detection of B_{80} by photoelectron spectroscopy. Density functional theory (DFT) computations revealed that the most stable structure includes a quasi-planar arrangement with two adjacent hexagonal holes. However, a cage structure is also viable as the relative energy to the global minimum is only 1.7 kcal·mol\(^{-1}\). In contrast, the lowest energy structure of the neutral B_{40} cluster possesses a perfect cage-like shape containing two hexagonal and four heptagonal holes (1, see Fig. 1) and it is more stable than the corresponding quasi-planar form by approximately 27 kcal·mol\(^{-1}\). The bonding network in B_{40} is completely delocalized via \(\sigma\)- and \(\pi\)-type multicenter bonds. The structure gains its stability due to its very high HOMO-LUMO energy gap (3.13 eV), which is comparable to that of C_{60} (3.02 eV). So, B_{40} represents the first pure boron buckyball (borospherene).

Very recently, a couple of boron buckyballs (B_{38} and B_{39}) were also reported in the literature. Four years ago, the B_{19} cluster was reported to have a perfect planar structure with a central filled pentagonal unit inside a B_{13} ring. Some of us subsequently found a remarkable fluxional behavior in the B_{19} cluster. B_{19} exhibits an almost free rotation of the internal pentagon-shaped hub within the co-planar B_{13} ring akin to a Wankel motor, which is evident from the Born-Oppenheimer Molecular Dynamics (BO-MD) simulations. In fact, the zero-point energy is sufficient to overcome the barrier that allows the inner and outer rings to rotate independent of each other. The presence of various multicenter bonds between the outer-ring and the inner-ring, that could easily migrate from one position to another during rotation, facilitates such a dynamical feature. This fluxional behavior is not limited to B_{19} only, B_{13}, B_{14}, B_{15}, B_{16}, B_{17} and B_{20} also show similar dynamical features. As B_{40} has a similar multicenter...
bonding pattern, there arises the compelling question, “Is B₄₀ also fluxional?” and that prompted us to explore the probable dynamical behavior of B₄₀. Indeed, a fluxional behavior of a cage would correspond to a nanobubble — a molten, hollow pure boron object, which is unprecedented to our knowledge, and thus is of immense interest concerning its chemical and physical properties.

In order to explore the dynamical behavior of B₄₀, we carried out a series of BO-MD simulations at the PBE/DZVP level in deMon2K (deMon2k v. 3, the deMon developers, Cinvestav, Mexico City 2011). The simulations are launched from the equilibrium geometry of B₄₀ (1) with random velocities assigned to the atoms, employing a Hoover thermal bath, for a simulation time of 25 ps with 0.5 fs time steps. During the MD computations, we keep the total angular momentum of the cluster as zero, thereby suppressing the cluster rotation. The behavior of the mean square displacement (msd) as a function of time easily allows us to differentiate between a solid-like and a liquid-like behavior. The mean-square displacement at time t is given by

\[
\langle \text{msd} \rangle = \frac{1}{N} \sum_{i=1}^{N} [r_i(t) - r_i(0)]^2 ,
\]

where \(r_i(t)\) is the position vector of the i-th atom at the time t and N is the total number of atoms in the system.

During the BO-MD simulations at 1000 K, the cluster maintains its connectivity pattern and cage-like structure as was reported by Zhai et al. But at 1200 K and 1500 K, a continuous transformation between the seven membered rings (7-MRs) and six membered rings (6-MRs) is perceived (see movie files in the electronic supporting information). The structure of the rear 7-MR first gets distorted by moving one B atom towards the adjacent 6-MR and then it becomes a 6-MR, which transforms the contiguous 6-MR into a 7-MR. But still the near 7-MR does not show any significant structural change. Thereafter, the rear 6-MR is again transformed into 7-MR and the near 7-MR starts to have a structural deformation. Therefore, the boron cage certainly shows a series of transformation in which the ring sizes continuously get changed during the simulations.

Figure 1 depicts the minimum energy structures and the transition state involved in the fluxional behavior of B₄₀. The relative energies (in kcal·mol⁻¹) are computed at the PBE0/6-311+G(d) level, including the zero point energy correction.

![Figure 1. The minimum energy structures (1 and 2) and the transition state (TS1) involved in the fluxional behavior of B₄₀. The relative energies (in kcal·mol⁻¹) are computed at the PBE0/6-311+G(d) level, including the zero point energy correction.](image-url)
surrounding the two 6-MRs and all of them are suitable for such a structural reorganization. If the central B₅ moiety changes, the original structure is recovered. In contrast, if any of the other B₅ fragments switches, then four boron atoms will change their positions to form a 7-MR from a 6-MR.

Real-temperature simulations on realistic time scales are computationally impossible for quantum-mechanical systems to date. As a typical time step is about one fs ($10^{-15}$ s), $10^{15}$ MD steps would be required. Given that it is unrealistic from the computational point of view, the standard procedure to overcome this problem is to run simulations at higher temperature, what essentially serves as a time lapse. In our case, the rare event is overcoming a barrier of 14.3 kcal/mol. Using Boltzmann’s law, our 25 ps simulation at 1200 K corresponds roughly to 0.5 ms in reality. Thus, with this – admittedly crude – approximation we do observe fluxionality at chemically relevant time scales. In other words, it does not mean that the transformation occurs at 1200 K, it is only that at this temperature, the molecule has the enough kinetic energy to cross the barrier in a 25 ps time frame.

The extraordinary dynamical behavior of B₄₀ is also reflected in the $\langle \text{msd} \rangle$ values (Fig. 2). In a solid-like system, one would expect the $\langle \text{msd} \rangle$ to be essentially a constant close to zero, whereas a liquid-like system should exhibit a nearly linear increase in $\langle \text{msd} \rangle$ with respect to time. At 1000 K, clearly B₄₀ behaves like a solid-like system, where the structure remains almost the same during the MD simulation, but the behavior of $\langle \text{msd} \rangle$ changes drastically at 1200 K and 1500 K. As a result of the nuclear mobility, B₄₀ can be described as a surface or two-dimensional liquid-like system at these temperatures. Thus, B₄₀ is a system that behaves like a support-free spherical two-dimensional liquid.

Fluxionality in carboranes is common, but no such interchange in the sizes of the rings occurs there. On the other hand, the classical carbon cages like C₆₀ do not show such type of transformations. A somewhat similar process is the formation of Stone-Wales (SW) defects in fullerenes or nanotubes. However, particularly for C₆₀, this isomerization is restricted due to a very high activation energy barrier (approximately 7 eV in fullerenes). So, the barrier in B₄₀ is significantly lower than that in the SW defects formation in carbon systems. In other words, B₄₀ is the first boron cluster that shows an interesting dynamical behavior with a moderate barrier (less than 15 kcal/mol), which allows that the ring sizes get changed continuously.

Figure 2. The profiles of the z-component of the induced magnetic field ($B^{\text{ind}z}$) for 1. The blue (red) line shows the magnetic response computed when the external magnetic field is applied perpendicular to the 6-MR (7-MR). The scale is given in ppm ($\mu$T for $|B^{\text{ext}}| = 1$ T).
The bonding situation in B₄₀ is intriguing. Zhai et al. studied in detail the nature of bonding in using the adaptive natural density partitioning (AdNDP) analysis and showed that a strong delocalization of the σ- and π-system is present. 48 pairs of valence electrons are delocalized via σ-bonds and the remaining electrons are delocalized via π-bonds. Zhai et al. also computed the nucleus independent chemical shift (NICS) to determine whether B₄₀ could be classified as an aromatic system. NICS indicates that the system is indeed aromatic in nature. However, there are some details about delocalization that were not explored previously and could provide some insight into the fluxional behavior of. Here we use the induced magnetic field, particularly the z-component of the induced magnetic field (B_{indz}), in order to understand delocalization.

Figure 3 shows the profiles of B_{indz} for external fields applied perpendicular to the 6-MRs. An external magnetic field in this direction can induce a current around the cage. The magnetic response at the cage center is very high (B_{indz} = -47.1 ppm) and diatropic in nature. This value is even larger than those computed in other spherenes. The intensity of B_{indz} in B₄₀ diminishes gradually along the center to surface of the cage. The radius of B₄₀ is approximately 2.5 Å. At this distance, the |B_{indz}| values are around 20 ppm lower than that computed at the cage center. However, at both 6- and 7-MRs, the magnitudes are still appreciable (B_{indz} ≈ -20 ppm), indicating a strong delocalization in the σ-framework.

The B_{indz} profile, which is mathematically equivalent to NICS_{zz}, shows a long-range shielding cone above the 6-MRs with an extension of nearly 10 Å, which is even more intense than that computed at C₆₀. In contrast, the magnetic response is different on the 7-MRs. In this case, the B_{indz} values become positive at around 3.5 Å above, implying a decrease in delocalization above the 7-MRs. Hence, these results indicate a strong π-delocalization at the 6-MRs, which is much higher than those at the 7-MRs.

The B_{indz} profiles of TS1 and 2 show essentially the same magnetic responses in both, shape and intensity as that in 1 (see Fig. 1-SI). Thus, during the rearrangement process, the changes in the σ- and π-delocalizations are not significant. This is the main difference compared to the carbon cages like C₆₀, where the changes in the connectivity (as Stone-Wales defects) modify the σ- and π-frameworks drastically and accordingly the delocalization and stability.

In summary, BO-MD simulations reveal that B₄₀, the first pure all-boron buckyball, shows a fascinating dynamical behavior in which the ring sizes continuously get changed during simulation. The transformation between 6- and 7-MRs occurs through an activation energy barrier of 14.3 kcal-mol⁻¹. The electronic structure of B₄₀ favors such transformation as the σ- and π-delocalizations are properly maintained in both the minimum energy structures (1 and 2) as well as in the corresponding transition.

Figure 3. Temporal variation of the mean square displacements, , for B40 at 1000, 1200, and 1500 K.
state (TS1). During the rearrangement, the changes around the multicenter B-B bonds are almost negligible, thus unlike in the carbon cages the barrier is not prohibitive. This is also related to the fact that a B-B bond is weaker than the C-C one. So, the dynamic behavior of $B_{40}$ is a consequence of the strong delocalization existing around the multicenter B-B bonds. The msd supports that $B_{40}$ is a nanobubble that behaves like a support-free spherical two-dimensional liquid.

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Author Contributions

G.M., J.L.C., A.D., and R.I. performed all the computations. S.P. and P.K.C. helped in analyzing the results and in writing the manuscript. T.H. and G.M. analyzed the results and wrote the manuscript. All authors reviewed the manuscript.

Additional Information

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