Generation of high energy laser-driven electron and proton sources with the 200 TW system VEGA 2 at the Centro de Laseres Pulsados

L. Volpe1,2, R. Fedosejevs3, G. Gatti1, J. A. Pérez-Hernández1, C. Méndez1, J. Apiñaniz1, X. Vaisseau1, C. Salgado1,4, M. Huault1,4, S. Malko1,4, G. Zeraouli1,4, V. Ospina1,4, A. Longman5, D. De Luis1, K. Li1, O. Varela1, E. García1, I. Hernández1, J. D. Pisonero1, J. García Ajates1, J. M. Alvarez1, C. García1, M. Rico1, D. Arana1, J. Hernández-Toro1, and L. Roso1,4

1Centro de Laseres Pulsados (CLPU), Edificio M5. Parque Cientifico. C/Adaja, 8. 37185 Villamayor, Salamanca, Spain
2Laser-Plasma Chair at the University of Salamanca, Salamanca, Spain
3Department of Electrical and Computer Engineering, University of Alberta, Edmonton, Alberta T6G 2V4, Canada
4University of Salamanca, Salamanca, Spain

(Received 21 December 2018; revised 17 January 2019; accepted 30 January 2019)

Abstract
The Centro de Laseres Pulsados in Salamanca, Spain has recently started operation phase and the first user access period on the 6 J 30 fs 200 TW system (VEGA 2) already started at the beginning of 2018. In this paper we report on two commissioning experiments recently performed on the VEGA 2 system in preparation for the user campaign. VEGA 2 system has been tested in different configurations depending on the focusing optics and targets used. One configuration (long focal length \( F = 130 \text{ cm} \)) is for underdense laser–matter interaction where VEGA 2 is focused onto a low density gas-jet generating electron beams (via laser wake field acceleration mechanism) with maximum energy up to 500 MeV and an X-ray betatron source with a 10 keV critical energy. A second configuration (short focal length \( F = 40 \text{ cm} \)) is for overdense laser–matter interaction where VEGA 2 is focused onto a 5 \( \mu \text{m} \) thick Al target generating a proton beam with a maximum energy of 10 MeV and temperature of 2.5 MeV. In this paper we present preliminary experimental results.

Keywords: high power laser; laser-plasma; particle acceleration

1. Introduction
Laser technology has advanced to the point where hitherto unobtainable intensities are now routinely achievable, and rapid progress is being made to increase intensities further[1]. The versatility of high power lasers has resulted in their use in a broad range of scientific fields, including novel particle accelerators, fusion research, laboratory astrophysics, condensed matter under high pressure, novel X-ray sources and strong-field quantum electrodynamics, among others. The potential for developing compact, high brightness particle and radiation sources has given a strong impetus to the development of the underpinning laser technology, including increasing the efficiency and repetition rate of the lasers. A result of this technological development can be seen in the new generation of ultrafast high power laser systems working at high repetition rate which have been built across Europe[2]. One of the most relevant and representative is the Centro de Laseres Pulsados[3] (CLPU) in Salamanca, Spain which has recently started operation. The CLPU has been founded by Spanish Ministry of Economy, Junta de Castilla y León and the University of Salamanca and its main system VEGA consists in a 30 fs pulse delivered in three different arms of 20 TW (VEGA 1), 200 TW (VEGA 2) and 1 PW (VEGA 3). CLPU has recently started operation phase, the first user access period on the VEGA 2 already started at the beginning of 2018 and a commissioning experiment on VEGA 3 is planned for 2019. VEGA 2 has been previously tested in different configurations depending on the focusing optics and targets used. One configuration is designed for underdense laser–matter interaction where VEGA 2 is focused \( (F = 130 \text{ cm}, \phi_L = 20 \mu \text{m}, Z_r = 260 \mu \text{m}) \) onto a low density gas-jet
generating (via wake field mechanism\cite{4–8}) electron beams with maximum energy up to 500 MeV and an X-ray betatron source with 10 keV characteristic critical energy\cite{9–12}. The second configuration is designed for over critical density laser–matter interaction where VEGA 2 is focused (\(F = 40\) cm, \(\phi_L = 7\) µm, \(Z_r = 25\) µm) onto a 5 µm Al target generating (via target normal sheath acceleration (TNSA) mechanism\cite{13–16}) a proton beam with a maximum energy of 10 MeV and average temperature of 2.5 MeV. In this paper both the commissioning experiments are reported and explained by describing the experimental setup and showing the capabilities of the VEGA system. Finally preliminary results are also given.

2. The VEGA system

The VEGA laser is a chirp pulse amplification (CPA) titanium:sapphire system, with a central wavelength of 800 nm ± 10 nm. It has three arms (see Figure 1) with maximum power 20 TW (VEGA 1), 200 TW (VEGA 2) and 1 PW (VEGA 3). The common front-end has a double CPA as well as an cross-polarized wave system that increases the contrast of the pulses significantly, as shown in Figure 2, making it suitable for high density targets.

The temporal contrast is: @ ns (pre-pulse) \(5 \times 10^{-10}\); 1 ps \(2 \times 10^{-5}\); 5 ps \(5 \times 10^{-8}\); @ 10 ps \(8 \times 10^{-9}\); 100 ps \(5 \times 10^{-12}\). Figure 2 shows a measurement of VEGA laser contrast just after the compressor.

In April 2017 the first call for user access on the VEGA 2 system was lunched by CLPU, and thirty scientific proposals were received for a total amount of more than 500 days requested over the 100 days offered in the call. The VEGA 2 system has been offered at high repetition rate with some limitations mainly connected to targetry and diagnostics. In preparation of such user access we organized a series of commissioning experimental campaigns along 2017. The main goal of this campaign was to learn how to prepare the VEGA 2 system for future user access campaigns. In particular we prepared the VEGA 2 target area to be easily adaptable for different possible experimental setups by changing the laser focusing system and by implementing a simple system to split the main beam (via beam splitters) at different fractions among which 1%/99% and 10%/90%, and other possibilities such as 50%/50% are under development. Depending on the experimental requirements a probe beam can be focused by one of the parabolic mirrors or simply by optical lenses. Both the parabolic mirrors are protected by thin fused silica pellicles placed just in front of them. Laser pulse duration is routinely measured close to the target position by using a second order auto-correlator and compared with second harmonic optimization measurements.

3. The commissioning experimental campaign

The VEGA 2 target area currently offers two different configurations with the possibility of several laser probes depending on the experimental requests. For both the configurations the laser VEGA 2 has been delivered at 30 fs, with a maximum power of 150 TW corresponding to a maximum energy at the entrance of the compressor of around 7 J. The conversion efficiency \(\eta\) from the entrance of the compressor up to the target is measured to be around 50% giving a total energy on target around 3.5 J. Assuming the same energy conversion efficiency for both the configurations the total laser intensity is estimated to be \(I_L[\text{W} \cdot \text{cm}^{-2}] \sim 5.8 \times 10^{21}/A[\text{µm}^2]\), where \(A[\text{µm}^2] = \pi r_L^2\).

3.1. Long focalization in gaseous targets

VEGA 2 was focused by an \(F/13\) parabolic mirror into a 5 mm thick gas jet placed at 30 cm far from the target chamber center (TCC), as shown in Figure 3. The laser beam waist has been measured (at low power) to be full width at half maximum (FWHM) 17.7 µm ± 0.5 µm with less than 50% of the total energy in the FWHM giving a peak

**Figure 1.** Schematic of the three VEGA systems.

**Figure 2.** Contrast of VEGA 2, measured by Sequioa.
intensity within the FWHM of around $1.1 \times 10^{19} \text{ W \cdot cm}^{-2}$. Pulse duration has been adjusted during the experimental campaign to find the best acceleration performance by considering the relation between plasma frequency and laser pulse duration. Fully ionized plasma density has been fixed to $1.2 \times 10^{19} \text{ cm}^{-3}$ (over one order of magnitude larger than the threshold for the bubble regime) to maximize betatron radiation emission by generating a broad electron beam energy spectrum and leading to an increased accelerated charges (see Figure 4).

Electron beams up to 500 MeV have been measured with an electron spectrometer composed by different magnet dipoles coupled with a lanex scintillating detector and an imaging system (see Figure 3). Betatron radiation has also been measured with typical synchrotron-like energy spectrum (see Figure 4) characterized by a critical energy of $\approx 10 \text{ keV}$. Betatron spectrum has been analysed with an X-ray CCD camera by using Ross filters technique. We estimated to have a peak brightness greater than $10^8 \text{ photons/sr at 0.1\% BW per shot on average, with a divergence greater than 9.2 \times 4.5 \text{ mrad in horizontal and vertical directions. This betatron source has been also used as high frequency (short wavelength) probe to investigate a pre-heated warm dense aluminium sample and experimental results are now under analysis and will be published soon.}

3.2. Short focalization in solid targets

VEGA 2 was focused by an $F/4$ gold coated parabolic mirror onto 5 $\mu$m thick Al foil with an angle of $10^\circ$ with respect to the normal of the foil (see Figure 5). The laser beam waist has been measured (at low power) to be FWHM 8 $\mu$m $\pm 2 \mu$m with 50% of the total energy in the FWHM giving a peak laser intensity within the FWHM of around $I_L \approx 2 \times 10^{20} \text{ W \cdot cm}^{-2}$. Laser pulse duration has been routinely measured to be 30 fs $\pm$ 3 fs. A three-axis motorized target holder has been used holding 5 and 10 $\mu$m thick Al foils. Target alignment was done prior to each shot with a resolution of around 5 $\mu$m ($\leq Z_r \sim 25 \mu$m).

For such laser intensities and so thin targets the control of laser pre-pulse\cite{17} and plasma corona expansion is mandatory. A first measurement of the laser contrast just after the compressor has been done by the laser team and it is shown in Figure 2. We developed a simple front-side diagnostic to control the laser–target interaction by looking at plasma formation as a function of time in a time range of $\pm 300 \text{ ps}$ with respect to the main beam arrival. A small part of the main beam is extracted, doubled in frequency, and its polarization rotated $90^\circ$ with respect to the main laser and focused (counterpropagating with respect to the main laser pulse) on the front target surface with an angle of $\approx 20^\circ$. The reflected beam was imaged onto an optical CCD camera, while a part of the transmitted beam is also recorded for reference. The setup is shown in Figure 5. Comparison of the amount of reflected probe beam at different delays will help for a preliminary interpretation and understanding of the laser contrast efficiency in laser–solid interaction; a further optimization of the diagnostic method will help to reconstruct the plasma density profile before, during and after the main pulse arrival by comparison to simulations. Figure 6 shows reflected fraction of the probe as a function of the chosen delay.

Proton beam energies up to 10 MeV have been routinely measured with different diagnostic techniques. Proton spatial and energy distributions have been measured first with
a stack of radio chromic films (RCFs) with the energy resolution limited by the number and thickness of RCF layers to $\Delta E \sim 0.5$ MeV. Time of flight (ToF) measurements have been performed by using 1 ns time-resolved pin diode detector and with a 300 ps time-resolved micro channel plate (MCP) both placed 2.5 m far from the TCC as shown in Figure 5. Proton measurements were also checked by CR39 detectors placed beside the RCF stack and the ToF detectors. The measurements show TNSA-like proton energy spectrum with an averaged (the average is done over different acquisitions) mean temperature $\langle T_{\text{mean}} \rangle \sim 2.5$ MeV, an averaged maximum temperature $\langle T_{\text{max}} \rangle \sim 9$ MeV and a divergence angle ranging between $20^\circ$ and $25^\circ$ FWHM. Figure 7 shows the RCF stack scheme with an example of experimental results.

Figure 8 shows proton energy spectra obtained by both RCFs and ToF analysis and in the insertion an example of ToF spectrum with maximum energy up to 10 MeV. Maximum proton energies from 6 to 10 MeV have been routinely obtained and an estimation of the total number of protons per shot ($\sim 10^8$) has been done by calibrating the RCF detectors with well known proton sources\cite{18}. ToF detectors are at the moment not absolutely calibrated and can only give information about the proton energy range.

Proton radiography of test objects has been performed with the obtained proton distribution and some results are shown in Figure 9 starting from left a piece of a leaf, a wing, a grid and the CLPU metal logo being radiographed. Analysis of the results is in progress and will be submitted soon for publication.

TNSA proton acceleration mechanism is controlled by electron dynamics into the target. It is relevant to have information on the laser-driven electron population\cite{19} which generates the quasi-static electric field\cite{20, 21} responsible for the proton acceleration. We imaged the electron beam travelling into the target by collecting and focusing incident K$_\alpha$ radiation coming from hot electron interaction within the Al target ($h\nu \sim 1.5$ keV). Such radiation has been routinely collected by a remotely adjustable Kirk-Patrick Baez (KB) microscope, onto an X-ray CCD. The KB microscope was placed 1.2 m from TCC at $30^\circ$ with respect to the target normal and the CCD camera placed in the KB microscope focal plane. This compact X-ray K$_\alpha$ diagnostic has also been tested by performing radiography of a calibration grid. Figure 10 shows the diagnostic setup, an example of an obtained K$_\alpha$ spot, and results from X-ray radiography of a grid pattern.

Proton production has been tested at a constant repetition rate mode and the VEGA 2 laser pulse operation has shown good stability in terms of shot-to-shot variation of
the focal spot and total laser energy which results in a good reproducibility of the experimental data in all the diagnostics. The $K_{\alpha}$ signal was used as a reference criterion for evaluating shot-to-shot fluctuation. Figure 11 shows an example of a series of shots comparing experimental data acquired from different installed diagnostics pin diode, multi-channel-plates and KB microscope. Clear correlation between results can be seen.

4. Conclusions

The first experimental campaigns on VEGA 2 target area have shown the potential of the 200 TW system demonstrating a good performance and stability of the VEGA 2 laser system with a large margin of improvement. A flexible and comprehensive targetry and a diagnostic system have been developed to support experimental activities in both long focal and short focal lengths configurations. It is worth noting that targetry and diagnostic systems need further development to match the repetition rate potential of the laser system. Indeed, target movement and alignment still need to be improved and optimized for working modes above 0.1 Hz. Diagnostic techniques are not always mature for high repetition rate acquisition that can reproduce the performance of the original passive detectors. In summary we have routinely obtained 0.5 GeV electron beams with more than 10 keV critical energy (synchrotron-like spectrum) X-ray betatron emission in the long focal length configuration and a TNSA-like proton spectrum with maximum energy of around 9 MeV in the short focal length configuration. In addition significant electromagnetic pulses have also been generated and measured during laser operation, in particular during the short focal campaign where laser intensities reach up to $10^{20}$ W·cm$^{-2}$.

Acknowledgements

The authors would like to thank all the CLPU workers that during the last 10 years were working hard in order to build and develop the facility, in particular, the CLPU scientific, laser and administrative divisions as well as the engineering and radio protection areas. Also, we would like to thank the Ministry of Science, the Castilla y Leon region and the University of Salamanca for the constant support during all the phases of construction and implementation of the CLPU facility. Support from Spanish Ministerio de Ciencia, Innovación y Universidades through the PALMA Grant No. FIS2016-81056-R, ICTS Equipment Grant No. EQC2018-005230-P; from LaserLab Europe IV Grant No. 654148 and from Junta de Castilla y León Grant No. CLP087U16 is acknowledged.

Author contributions

R. Fedosejevs and L. Volpe conceived, designed and oversaw respectively the electron acceleration and the proton acceleration experiments; G. Gatti and J. A. Perez supervised and operated both the experiments; X. Vaisseau, J. I. Apiñaniz, M. Huault, C. Salgado, G. Zeraouli, S. Malko, A. Longman, K. Li, and V. Ospina, performed the experiments and analysed the data; C. Mendez and Laser Team operated the laser and D. de Luis with engineer area supported experimental campaign with design and fabrication. L. Volpe wrote the first draft of the paper.

References

1. D. Strickland and G. Mourou, Opt. Commun. 56, 219 (1985).
2. C. Danson, D. Hillier, N. Hopps, and D. Neely, High Power Laser Sci. Eng. 3, e3 (2015).
3. www.clpu.es.
4. T. Tajima and J. M. Dawson, Phys. Rev. Lett. 43, 267 (1979).
5. S. P. D. Mangles, C. D. Murphy, Z. Najmudin, A. G. R. Thomas, J. L. Collier, A. E. Dangor, E. J. Divall, P. S. Foster, J. G. Gallacher, C. J. Hooker, D. A. Jaroszynski, A. J. Langley, W. B. Mori, P. A. Norreys, F. S. Tsung, R. Viskup, B. R. Walton, and K. Krushelnick, Nature 431, 535 (2004).

6. J. Faure, Y. Glinec, A. Pukhov, S. Kiselev, S. Gordenko, E. Lefebvre, J.-P. Rousseau, F. Burgy, and V. Malka, Nature 431, 541 (2004).

7. C. G. R. Geddes, C. Toth, J. Van Tilborg, E. Esarey, C. B. Schroeder, D. Bruhwiler, C. Nieter, J. Cary, and W. P. Leemans, Nature 431, 538 (2004).

8. V. Malka, Phys. Plasmas 19, 055501 (2012).

9. F. Albert, R. Shah, K. T. Phuoc, R. Fitour, F. Burgy, J.-P. Rousseau, A. Tafzi, D. Douillet, T. Lefrou, and A. Rousse, Phys. Rev. E 77, 056402 (2008).

10. S. Kneip, C. McGuffey, J. L. Martins, S. F. Martins, C. Bellei, V. Chvykov, F. Dollar, R. Fonseca, C. Huntington, G. Kalintchenko, A. Maksimchuk, S. P. D. Mangles, T. Matsuoka, S. R. Nagel, C. A. J. Palmer, J. Schreiber, K. Ta Phuoc, A. G. R. Thomas, V. Yanovsky, L. O. Silva, K. Krushelnick, and Z. Najmudin, Nat. Phys. 6, 980 (2010).

11. M. Z. Mo, Z. Chen, S. Fourmaux, A. Otani, J. C. Kieffer, Y. Y. Tsui, A. Ng, and R. Fedosejevs, Rev. Sci. Instrum. 84, 123106 (2013).

12. F. Albert, Plasma Phys. Control. Fusion 58, 103001 (2016).

13. S. C. Wilks, A. B. Langdon, T. E. Cowan, M. Roth, M. Singh, S. Hatchett, M. H. Key, D. Pennington, A. MacKinnon, and R. A. Snively, Phys. Plasmas 8, 542 (2001).

14. A. Maksimchuk, S. Gu, K. Flippo, D. Umstadter, and V. Y. Bychenkov, Phys. Rev. Lett. 84, 4108 (2000).

15. E. L. Clark, K. Krushelnick, J. R. Davies, M. Zepf, M. Tatarakis, F. N. Beg, A. Machacek, P. A. Norreys, M. I. K. Santala, I. Watts, and A. E. Dangor, Phys. Rev. Lett. 84, 670 (2000).

16. R. A. Snively, M. H. Key, S. P. Hatchett, T. E. Cowan, M. Roth, T. W. Phillips, M. A. Stoyer, E. A. Henry, T. C. Sangster, M. S. Singh, S. C. Wilks, A. MacKinnon, A. Offenberger, D. M. Pennington, K. Yasuake, A. B. Langdon, B. F. Lasinski, J. Johnson, M. D.erry, and E. M. Campbell, Phys. Rev. Lett. 85, 2945 (2000).

17. M. Kaluza, J. Schreiber, M. I. K. Santala, G. D. Tsakiris, K. Eidmann, J. Meyer ter Vehn, and K. J. Witte, Phys. Rev. Lett. 93, 045003 (2004).

18. Centro de Micro-Análisis de Materiales Universidad Autónoma de Madrid, http://www.cmam.uam.es/es/instalaciones/lineas-de-haz/79.

19. L. Volpe, D. Batani, G. Birindelli, A. Morace, P. Carpegiani, M. H. Xu, F. Liu, Y. Zhang, Z. Zhang, X. X. Lin, F. Liu, S. J. Wang, P. F. Zhu, L. M. Meng, Z. H. Wang, Y. T. Li, Z. M. Sheng, Z. Y. Wei, J. Zhang, J. J. Santos, and C. Spindloe, Phys. Plasmas 20, 033105 (2013).

20. J. Schreiber, F. Bell, F. Gruner, U. Schramm, M. Geissler, M. Schnürer, S. Ter-Avetisyan, B. M. Hegelich, J. Cobble, E. Brambrink, J. Fuchs, P. Audebert, and D. Habs, Phys. Rev. Lett. 97, 045005 (2006).

21. M. Passoni and M. Lontano, Phys. Rev. Lett. 101, 115001 (2008).

22. L. Volpe, D. Batani, B. Vauzour, Ph. Nicolai, J. J. Santos, C. Regan, A. Morace, F. Dorcieres, C. Fourment, S. Hulin, P. S. Baton, K. Lancaster, M. Galimbet, R. Heathcote, M. Tolley, C. Spindloe, P. Koester, L. Labate, L. A. Gizzi, C. Benedetti, A. Sgattoni, M. Richetta, J. Pasley, F. Beg, S. Chawla, D. P. Higginson, and A. G. MacPhee, Phys. Plasmas 18, 006101 (2011).

23. M. Z. Mo, Z. Chen, S. Fourmaux, A. Saraf, K. Otani, R. Masoud, J.-C. Kieffer, Y. Tsui, A. Ng, and R. Fedosejevs, Phys. Rev. E 95, 053208 (2017).