First Simulation Results for the Photon in a Non-Commutative Space

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We present preliminary simulation results for QED in a non-commutative 4d space-time, which is discretized to a fuzzy lattice. Its numerical treatment becomes feasible after its mapping onto a dimensionally reduced twisted Eguchi-Kawai matrix model. In this formulation we investigate the Wilson loops and in particular the Creutz ratios. This is an ongoing project which aims at non-perturbative predictions for the photon, which can be confronted with phenomenology in order to verify the possible existence of non-commutativity in nature.

1. THE NON-COMMUTATIVE U(1) MODEL

We investigate non-perturbatively QED on a non-commutative (NC) space-time [1], where a commutation relation of the form

\[ [\hat{x}_\mu, \hat{x}_\nu] = i\Theta_{\mu\nu} \]

is introduced. This has the effect of replacing points in regular space with “Planck cells”. \( \Theta_{\mu\nu} \) is a real and antisymmetric non-commutativity tensor.

Here we address pure U(1) gauge theory in a Euclidean NC space-time. Our main tool is the star product, which is defined as

\[ f(x) \ast g(x) = f(x) \exp\left(\frac{i\Theta_{\mu\nu}}{2} \partial_\mu \partial_\nu \right) g(x). \]

This definition absorbs the “space deformation” into the products of fields, thereby allowing us to treat the space-time operators \( \hat{x} \) and functions thereof as regular space-time coordinates again,

\[ \hat{\phi}(\hat{x}) \hat{\psi}(\hat{x}) \rightarrow \phi(x) \ast \psi(x). \]

Equipped with this tool we are able to write down the gauge action of the pure U(1) theory as

\[ S[A] = -\frac{1}{4g^2} \int d^4x \text{Tr}[F_{\mu\nu}(x) \ast F_{\mu\nu}(x)], \]

where \( F_{\mu\nu} \) is given by

\[ F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu - i(A_\mu \ast A_\nu - A_\nu \ast A_\mu). \]

The gauge fields in this theory do not commute, as the Yang-Mills term shows. Another unusual feature of this theory are complex, \( \ast \)-gauge invariant Wilson loops [2]. Moreover a \( \Theta \)-deformed dispersion relation for the photon is expected [3]. Such a modification was observed for the NC \( \lambda\phi^4 \) model [4]. This is due to a remarkable feature of the NC field theories, the UV/IR mixing of divergences [5]. That property makes the perturbative treatment beyond one loop extremely difficult.

2. LATTICE FORMULATION OF THE NC PLANE

We carried out simulations involving two NC directions, as well as two commutative directions, which include the Euclidean time. To be explicit, we set \( \Theta_{12} = -\Theta_{21} = \theta \), and \( \Theta_{\mu\nu} = 0 \) otherwise. The parameter \( \theta \) was assumed to be constant throughout the space-time.

In order to perform Monte Carlo simulations one needs to formulate the theory on a lattice. This is achieved by the operator relation [6]

\[ \exp\left(\frac{2\pi i}{a} \hat{\xi}_\mu \right) = \hat{1}, \quad \mu = 1, 2, \]

which imposes a discrete non-commutative space-time. After having obtained the lattice space-time, it needs to be populated by U(1) link variables. The key to render their simulation — involving the non-local \( \ast \)-product — feasible is the
Morita equivalence \([7]\) of NC U(\(N\)) gauge theories to the twisted Eguchi-Kawai (TEK) model \([8]\), with the action
\[
S_{TEK}[U] = -\beta \sum_{\mu \neq \nu} Z_{\mu \nu} \text{Tr} \left( U_\mu U_\nu U_\nu^\dagger U_\mu^\dagger \right). \quad (7)
\]

Here the \(U_\mu\) are U(\(N\)) matrices. The twist \(Z_{\mu \nu} = Z_{\nu \mu}^*\) is given by \(\exp\left(\frac{2\pi i n_{\mu \nu}}{N}\right)\), with \(n_{\mu \nu} \in \mathbb{Z}\). In our mapping we chose \(n_{21} = -n_{12} = (N + 1)/2\) for some odd \(N\), and \(n_{\mu \nu} = 0\) otherwise.

We discretized the commutative plane to a regular \((N-1) \times (N-1)\) lattice, where on each lattice site the NC plane is represented by a \(N \times N\) TEK model (\(N\) is odd).

The non-commutativity parameter \(\theta\) and the lattice parameters \(N\) and \(a\) are related as \([8]\)
\[
\theta = \frac{1}{\pi} N a^2. \quad (8)
\]

In the \(N \to \infty, a \to 0\) limit this theory describes a continuum gauge theory, which theory one obtains depends on the way this limit is taken. Keeping \(\theta\) finite in the continuum and infinite volume corresponds to the double scaling limit. In the case of finite \(N\) and \(a\) the theory maps exactly to a NC U(1) lattice theory \([9]\).

In order to be able to apply the heat bath algorithm, the action was linearized in \(U_\mu\), as a generalization of the scheme of Ref. \([9]\).

3. RESULTS

In the simulation results presented here we set \(N = 25\). The upper part of Fig. 1 shows the real part of the plaquette values at different \(\beta\). In this case all six possible orientations of the plaquettes are averaged. The plot shows that the results match the anticipated asymptotic behavior \([8]\). The phase transition occurs at around \(\beta \approx 0.35\).

The decay of Wilson loops observed from differently sized squares is shown in the lower part of Fig. 1. In that plot the results are split according to the plaquette orientations (NC, mixed and commutative). The real part of the NC Wilson loops oscillates, while eventually decaying towards zero. This is in qualitative agreement with the observations in Ref. \([10]\) for \(d = 2\). For the other two types we see a monotonous decay towards zero. Notice that these results were obtained in the weak coupling phase.

We have calculated Creutz ratios \(\chi\) for various rectangular Wilson loops. The results are shown in Fig. 2. From these we see that the string tension seems to approach zero as \(\beta \to \infty\), which corresponds to the continuum limit.

Fig. 3 shows the correlation function of the plaquettes lying completely in the NC plane, and separated by \(\Delta t\) in Euclidean time. These data were taken at \(\beta = 0.4\), so we are barely in the weak coupling phase. The decay seems to be exponential, but more statistics are required to confirm this behavior.

Fig. 4 shows the commutative analog to the correlation function discussed above. However, in that case the decay does not seem exponential.
4. SUMMARY AND CONCLUSION

We have simulated a NC U(1) field theory. First observables such as the action, plaquettes, Wilson loops and Creutz ratios could be measured. The results for the action and plaquettes agree with asymptotic predictions. There seems to be a vanishing string tension at weak coupling in the part of the parameter space we explored.

Further work will hopefully allow us to identify a physical scale and to extrapolate to a NC continuum limit by means of double scaling (i.e. a fixed product $N_a^2$). Thus we hope to glimpse at the dispersion relations for the NC photon, which might allow us to confront the theory with nature.

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