Beam dynamics performances and applications of a low-energy electron-beam magnetic bunch compressor

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Abstract

Many front-end applications of electron linear accelerators rely on the production of temporally-compressed bunches. The shortening of electron bunches is often realized with magnetic bunch compressors located in high-energy sections of accelerators. Magnetic compression is subject to collective effects including space charge and self interaction via coherent synchrotron radiation. In this paper we explore the application of magnetic compression to low-energy (∼40 MeV), high-charge (nC) electron bunches with low normalized transverse emittances (<5 µm).

Keywords: Photoinjector, Linear accelerator, Electron beam, Magnetic bunch compressor, Space charge, Coherent synchrotron radiation, Flat beam

1. Introduction

Most of the photoinjectors being used for generation of bright electron bunches for, e.g., free-electron laser (FEL) applications consist of generating and rapidly accelerating the electron bunch to high energy and subsequently shortening the bunch using a magnetic bunch compressor \cite{1}. The only deviation to such a design is the combined acceleration and compression using velocity bunching \cite{2}. Attempts to operate low-energy magnetic bunch compressors have to date been inconclusive \cite{3} or deemed incompatible with the production of low-emittance beams \cite{4}. In this study we explore and present via numerical simulations that low-energy bunch compression performed on a ∼40-MeV electron beam can be viable depending on requirements. We especially present trade-off curves between transverse emittance and peak current for several cases of electron-bunch charge. In addition, our simulations are performed with several computer programs thereby enabling a benchmarking of very different approaches for modeling collective effects and especially coherent synchrotron radiation (CSR) \cite{5, 6}. Our study considers the magnetic bunch compression planned in the 40-50 MeV photoinjector of the Advanced Superconducting Test Accelerator (ASTA) currently under construction at Fermilab \cite{7, 8}.

2. Accelerator beamline overview

The compression of a ∼40-MeV electron bunch via magnetic compression is investigated for the case of the ASTA photoinjector diagrammed in Fig. 1. The beamline includes a photoemission electron source consisting of a cesium telluride (Cs\textsubscript{2}Te) photocathode located on the back plate of a 1+1/2 cell radiofrequency (RF) cavity operating at 1.3 GHz \cite{9}. The cathode is illuminated with a 3-ps ultraviolet laser pulse with uniform radial distribution and a Gaussian temporal profile. The RF gun is surrounded by two solenoidal lenses that control the beam’s transverse size and emittance. Downstream of the RF gun, the typical beam energy is ∼5 MeV. The
bunches are further accelerated up to 50 MeV by two 1.3-GHz superconducting RF (SCRF) accelerating cavities (labeled as CAV1 and CAV2 in Fig. 1). A third SCRF cavity (CAV39) operating at 3.9 GHz will eventually be incorporated to correct for nonlinear longitudinal phase space distortions. Because of its superconducting nature, the ASTA facility produces electron bunches repeated at 3 MHz arranged in a 1-ms 5-Hz RF macropulse. The downstream beamline includes quadrupoles, steering dipole magnets, and diagnostics stations. A skew-quadrupole channel can be set up as a round-to-flat-beam transformer (RFBT) to convert an incoming angular-momentum-dominated beam into a flat beam with high transverse emittance ratio. The beamline also incorporates a four-bend magnetic bunch compressor (BC1) which, consists of four 0.2-m rectangular dipoles (B1, B2, B3, B4) with respective bending angles of (+,-,-,+) 18°. The longitudinal dispersion of BC1 is $R_{66} = -0.19$ m. Finally a single-shot longitudinal phase space diagnostics combining a transverse-deflecting cavity (TDC) with a vertical spectrometer will be installed.

The beam dynamics through CAV2 were simulated with astra and optimized using a genetic optimizer for several cases of charge and photocathode drive-laser configurations; see Ref. [16]. The resulting phase space distributions are used as a starting point for transport and compression through the beamline downstream of CAV39. The quadrupoles settings were optimized for the various operating charges using the single-particle dynamics program elegant. The evolution of the nominal betatron functions downstream of CAV39 up to the cryomodule entrance is plotted in Fig. 2.

![Figure 1: Injector configuration at ASTA. The “RF gun”, “L1” and “L2” respectively correspond to the gun cavity and surrounding solenoid magnets, “CAV1”, “CAV2”, and “CAV39” are superconducting RF cavities, “RFBT” is the round-to-flat beam transformer, and “BC1” refers to the magnetic bunch compressor, and B1-4 are the dipoles of the chicane, with distance between the dipoles marked in the figure. The number below the beamline indicates the axial positions in meters w.r.t. the photocathode surface.](image1)

![Figure 2: Evolution of the horizontal (blue) and vertical (red) betatron functions through the ASTA injector. The green rectangles indicate the location of quadrupole and dipole (smaller rectangles) magnets. The BC1 compressor is located at $s \in [11.9, 15.1]$ m. The origin of the horizontal axis ($s = 0$ m, not shown) corresponds to the photocathode surface. (color online)](image2)

3. Modeling methodology

The bunch-compression performance of BC1 was explored via numerical simulations using impact-z and csrtrack. Both programs model the beam as an ensemble of interacting macroparticles and integrate the equations of motion to advance the macroparticles along a user-specified beamline.

In impact-z the space-charge (SC) interaction is modeled using a mean-field quasi-static particle-in-cell (PIC) algorithm and the (point-like) macroparticles are advanced through the beamline using high-order transfer maps. Each beamline element is segmented into axial slices modeled by transfer maps. Between each transfermap segment, impact-z applies a space-charge “kick” evaluated from the mean-field PIC SC algorithm. CSR effects are included in impact-z using the one-dimensional formulation described in Ref. [20]. The one-
dimensional model is valid provided \[ \frac{\sigma_s(s)}{\sigma_z(s)} \ll 1, \quad \text{with} \quad \mathcal{D}(s) = \frac{\sigma_s(s)}{\sigma_z(s)} \left( \frac{\sigma_z(s)}{R(s)} \right), \] (1)

where \( R(s) \) is the trajectory’s radius of curvature and \( \sigma_s(s) \) and \( \sigma_z(s) \) are respectively the transverse and longitudinal root-mean-square (RMS) sizes at the curvilinear beamline position \( s \).

In \textsc{impact-z} the longitudinal charge distribution needed for the 1D CSR model is obtained from a longitudinal binning of the macroparticle ensemble. Convergence studies were carried out in order to determine the optimal number of longitudinal bins, \( N_z \), to be used for both the SC and CSR calculations. Low values of \( N_z \) generally underestimated the peak-current and therefore the collective effects, while large values of \( N_z \) introduce numerical noise that can lead to artifacts (e.g. numerically-induced microbunching) \[ [22]. \] The convergence study \[ [23] \] revealed an appropriate value of \( N_z = 256 \) for a bunch represented by \( N = 2 \times 10^5 \) macroparticles. The number of bins in the transverse dimensions was set to \( N_t = N_y = 16 \). The \textsc{impact-z} simulations presented in the rest of this paper use this set of parameters.

The program \textsc{csrtrack} was specifically developed to efficiently simulate the impact of bunch radiative self-interaction via CSR. \textsc{csrtrack} incorporates several models including a 2D particle-to-particle (P2P) model that directly computes the forces on macroparticles from the Liénard-Wiechart potentials evaluated at retarded times. These calculations are self-consistent and enable the computation of both the transverse and longitudinal force contributions from SC and CSR effects. Since the P2P model is computationally intensive, \textsc{csrtrack} also includes an improved one-dimensional model referred to as the 1D Projected (1DP) model. The 1DP model uses the 1D projection of the smoothed charge distribution convoluted with a kernel function \[ [24]. \] Compared to the model of Ref. \[ [21], \] the 1DP model is not limited to the ultra-relativistic regime. \textsc{csrtrack’s} P2P model treats each macroparticle as a 3D Gaussian charge distribution (referred to as “sub-bunch”) in the \((x,y,z)\) space with distribution

\[
g(x,y,z) = \frac{1}{(2\pi)^{3/2}\sigma_x\sigma_y\sigma_z} e^{-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2} - \frac{z^2}{2\sigma_z^2}}, \quad (2)
\]

where \( \sigma_x, \sigma_y \) and \( \sigma_z \) are respectively the horizontal, vertical and longitudinal RMS sizes of the sub-bunches. The resulting beam’s spatial charge distribution is \( \Phi(x,y,z) = \sum_{j=1}^{N} Q_j g(x-x_j,y-y_j,z-z_j) \) where \( Q_j \) and \((x_j,y_j,z_j)\) are respectively the \( i\)-th sub-bunch charge and LPS coordinates. In \textsc{csrtrack}, \( \sigma_x \) and \( \sigma_y \) may be defined relative to the vertical and longitudinal RMS bunch sizes, \( \sigma_y \) and \( \sigma_z \), respectively, and are adjusted along the bunch compressor as the dimensions of the bunch change. Due to the computational intensiveness of the P2P model, only \( 10^4 \) sub-bunches were used compared to \( 2 \times 10^5 \) used with the 1DP model. This relatively-low number of sub-bunches requires a large \( \sigma_y \) of 5\% of the longitudinal bunch length \( \sigma_z \). For the 1DP simulations and 10\% for the P2P simulations; for a detailed study see Ref. \[ [23]. \] The P2P model requires parameters for the sub-bunches in the horizontal and vertical dimensions, \( \sigma_y \) and \( \sigma_z \). We chose \( \sigma_y = 0.10\sigma_y \) and \( \sigma_z = 0.10\sigma_z \). However, \textsc{csrtrack} does not allow \( \sigma_y \) to be set as a variable of the horizontal RMS width, so we instead chose 0.1 mm, which is on the order of 10\% of the RMS size \( \sigma_y \) for all four of the bunch charges presented here. Both of \textsc{csrtrack’s} models neglect collective forces in the vertical dimension.

The \textsc{asta} simulations of the beam generation and acceleration simulated up to \( s = 8.2 \) m from the photocathode were used as a starting point for our simulations. A \textsc{python program, gluetrack} \[ [25], \] was used to manipulate the beam distributions and generate macroparticle distributions suitable for \textsc{impact-z} and \textsc{csrtrack}. For all the bunch-compressor studies, \textsc{impact-z} was used to track the bunch distribution from \( s = 8.2 \) up to \( s = 11.8 \) m corresponding to 0.1 m upstream of the entrance face of the first dipole (B1) of BC1. This distribution was used as a starting point for the bunch-compressor simulations, which include BC1 and a 1.0-m downstream drift to allow for the influence of SC possible transient CSR effects.

Simulations were performed for four cases of bunch charges ranging from 3.2 nC to 20 pC. For each charge the transverse emittance was optimized with \textsc{asta} \[ [16]. \] The distributions were manipulated using \textsc{gluetrack} to adjust several parameters including their Courant-Snyder (C-S) parameters, and longitudinal phase space (LPS) chirp \( C \equiv -(z_i\delta_i)/\sigma_{zi}^2 \) where \((z_i,\delta_i)\) are the coordinates in the LPS, the \( \langle u \rangle \) indicates the statistical averaging of variable \( u \) over the LPS distribution and \( \sigma_{zi} \equiv \langle z_i^2 \rangle^{1/2} \). We
also modeled the effect of CAV39 by numerically removing the second order correlation in the LPS distribution.

Table 1: Transverse and longitudinal beam parameters 0.1-m upstream of B1 dipole entrance face. Only the Courant-Snyder parameters were fixed while the other parameters depend on the bunch charge or upstream beamline settings.

| Parameter     | Value | Units |
|---------------|-------|-------|
| $\beta_{x,i}$ | 8     | m     |
| $\alpha_{x,i}$| 3     | -     |
| $\beta_{y,i}$ | 1.6   | m     |
| $\alpha_{y,i}$| -1.6  | -     |
| $C$           | 1.0, 6.0 | m$^{-1}$ |
| total energy  |       | 38.6 MeV |

The charge-independent beam parameters computed 0.1-m upstream of dipole magnet B1 are summarized in Tab. 1. The beam distribution was matched to achieve the transverse C-S parameters listed in Tab. 1 (see discussion below) upstream of BC1. The initial LPS distribution was tuned by removing the second-order correlation and scaling the first-order correlation between $\delta_i$ and $z_i$. The other LPS parameters and emittances (shown for the four charges in Tab. 2) are inherent to the generation process and were not adjusted. The initial LPS distribution for each of the four charges appears in Fig. 3 with its linear correlation removed. The $S$-shaped LPS is a remnant of space charge effects during the bunch generation and transport before acceleration in CAV1 and CAV2 [26]. As expected larger charges yield higher total fractional momentum spread.

Table 2: Initial normalized transverse $\varepsilon_{x,i}$ and longitudinal $\varepsilon_{z,i}$ emittances and RMS bunch length $\sigma_{z,i}$ for the four cases of charge considered in this paper. The parameters are computed 0.1-m upstream of dipole magnet B1’s entrance face.

| Q (nC) | $\varepsilon_{x,i}$ (µm) | $\varepsilon_{y,i}$ (µm) | $\varepsilon_{z,i}$ (µm) | $\sigma_{z,i}$ (mm) |
|--------|---------------------------|--------------------------|-------------------------|---------------------|
| 3.2    | 4.43                      | 4.58                     | 82.19                   | 2.56                |
| 1.0    | 2.20                      | 2.22                     | 33.41                   | 1.95                |
| 0.250  | 0.580                     | 0.576                    | 14.37                   | 1.93                |
| 0.020  | 0.296                     | 0.297                    | 2.54                    | 1.26                |

Simulations were performed for LPS chirps $C \in [1.0, 6.0] \text{ m}^{-1}$ with the bulk of the simulations performed around the maximum-compression value $C = -1/R_{36} \approx 5.2 \text{ m}^{-1}$, corresponding to enhanced collective effects.

The initial values for the betatron functions were selected such that the beam experiences a waist between the third and fourth dipole [27]. Simulations performed with the 1DP CSRTRACK model also confirmed there is a set of horizontal C-S parameters that minimizes the bending-plane emittance growth as displayed in Fig. 4. The upstream magnets were tuned to provide the incoming horizontal C-S parameter $(\beta_{x,i}, \alpha_{x,i}) = (8.0 \text{ m}, 3.0)$ shown in Fig. 2. Lastly, the final RMS bunch lengths as a function of the initial LPS chirp are given in Fig. 5 as obtained using IMPACT-Z’s SC+CSR model. In addition, there are slight variations in the bunch length between the four different models due to their collective effects inhibiting compression to varying degrees.

### 4. Benchmarking of numerical models

The beam dynamics simulations throughout BC1 were performed for several degrees of bunch compression (controlled with the LPS chirp) for the four cases of bunch...
charges that will be used at the ASTA. Four simulation algorithms were used as described earlier. Because of its computational effort, crtrack’s P2P model simulations were restrained to a smaller range of values for $C$.

The LPS 1.0-m downstream of the B4 dipole simulated with impact-Z (SC + one-dimensional CSR models) and crtrack (P2P model) are shown in Fig. 6 [(a)-(d)] and [(e)-(f)] respectively for the four charges listed in Tab. 2. Despite the vastly different algorithms used by these two programs, the LPS distributions displayed very similar distortions including those at the small-scale levels. Figure 7 summarizes the evolution of peak current as a function of the initial LPS chirp for the four numerical models.

Table 3: Final normalized longitudinal $\epsilon_z$ at maximum compression ($C = 5.2$ m$^{-1}$) simulated with impact-z (SC+CSR) and crtrack (P2P).

| Q (nC) | $\epsilon_z$ (\(\mu\)m) | crtrack (P2P) |
|--------|-----------------|---------------|
| 3.2    | 267             | 261           |
| 1.0    | 118             | 105           |
| 0.250  | 61.5            | 57.8          |
| 0.020  | 10.5            | 11.6          |

SC+CSR model, as these are the only two models that account for both SC and CSR effects. However, the P2P model also includes transverse CSR forces and a more elaborate model for longitudinal CSR. Our previous study [23] showed that the influences on final emittance from using too-few macroparticles as well from the randomization in the down-sampling of initial distributions were both much smaller than the discrepancy
between the SC+CSR and P2P models. The emittance growth observed from \textsc{csrtrack}’s 1D and \textsc{impact-z}’s SC-only model, indicates that CSR accounts for most of the emittance dilution at higher charge. For the low-charge simulations ($Q = 250$ and 20 pC) the relative importance is reversed, with SC contributing more to the emittance degradation than CSR. Of the models presented here, only \textsc{impact-z}’s include SC in both the vertical and the horizontal planes, and are shown in Fig. 9. As vertical emittance growth is entirely the result of SC, the inclusion of \textsc{impact-z}’s CSR model reduces vertical emittance growth due to the reduced compression.

Finally, the evolution of the slice parameters during compression were explored. For this analysis the beam is divided into axial slices of equal longitudinal length $\delta z = 20 \, \mu m$. A statistical analysis on the population contributing to each slice was performed to yield the slice emittances, energy spread and peak current. A comparison of the slice bending-plane emittance and energy spread between \textsc{impact-z} and \textsc{csrtrack}’s P2P model of slice-emittance as the chirp is varied appears in Fig. 10. The level of agreement between the two programs is of the same order as what observed for the bunch parameters. Fig. 11 summarizes the evolution of the slice horizontal emittance in the slice with the highest peak current within the bunch. When the beam is greatly over-compressed (C = 6.0 m$^{-1}$), the LPS may be double-peaked, and the transverse brightness as defined here may not be an appropriate measure of the bunch’s utility. The curve demonstrates how little slice emittance growth occurs for partial
5. Expected beam dilution and trade-offs

A large number of accelerator applications require beams with high-peak-currents and low-transverse-emittances. These requirements conflict with each other as collective effects, which dilute the beam’s phase space and emittances, increase with peak current. A commonly-used figure of merit is peak transverse brightness $B_\perp \equiv \frac{I}{\pi \varepsilon_x \varepsilon_y}$. Figure 12 summarized the evolution of $B_\perp$ as function of the LPS chirp for the four cases of bunch charges. The figure combines the data provided in Fig. 7 and Fig. 8. Despite the lower-charge bunch result in smaller peak current at lower charges (see Fig. 7), the transverse brightness increases with lower bunch charges. The main factor at play in this reduction is the lower initial transverse emittances, and more importantly, the reduced dilution of the transverse emittances during compression in BC1 due to the weaker collective effects (CSR and SC).

In addition we note that the maximum achieved transverse brightness does not necessarily occur at maximum compression. This is due to the larger peak currents at maximum compression that drive collective effects, wherein the relative emittance growth driven by collective effects is greater than the relative increase in peak current, a trade-off similar to that of going down to lower bunch charges. A summary of the achieved maximum value of $B_\perp$ appears in Fig. 13.

The trade-off between obtained peak current and $\varepsilon_x$ is shown in Fig. 14 only data associated to LPS chirp up to maximum compression at $C \leq 5.2 \text{ m}^{-1}$ are displayed, as over-compression results in lower peak currents with generally larger emittance dilutions.

The simulation data points to several conclusions about the parametric trade-offs that must be considered. First, the emittance growth, particularly the slice emittance, is greatly reduced at lower degrees of compression, particularly $C < 4 \text{ m}^{-1}$, which corresponds to compression
Figure 10: Example of final normalized transverse slice emittances (top row) and energy spread (bottom row) evolution within a 1-nC bunch for four cases of compression \( C = 4.0, 5.0, 5.25, \) and \( 6.0 \text{ m}^{-1} \) respectively shown as blue, green, red and turquoise traces. Plots (a) and (c) correspond to \texttt{impact-z} simulations while plots (b) and (d) are results from \texttt{csrtrack}'s P2P model. Emittance and energy spread values associated to slices that contain too-few number of macroparticles for meaningful statistical analysis are set to zero. The heads and tails of the bunches are sparsely populated (see Fig. 6 for reference), particularly for the P2P simulations which use only 5 % of the number of particles used in the \texttt{impact-z} simulations.

It is anticipated that a second-stage bunch compressor will eventually be installed at higher energy to compress the bunch to high peak current while mitigating phase-space dilution. However, in light of the studies presented in the previous Sections it is worth investigating other possible applications of the BC1 compressor as discussed below.

6.1. Multi-stage bunch compression

One of the motivations for exploring low-energy bunch compression is to moderately compress the bunches before injection in an accelerating structure and then further compress at high energy \[4\]. In Ref. \[4\], the low-energy bunch compression was accomplished at \(- 15 \text{ MeV} \) and resulted in intolerable beam degradation at the Tesla Test Facility I (TTF1) and consequently dismantled from the beamline during its upgrade as the Free electron LASer in Hamburg (FLASH) user facility. The requirement on the first-stage bunch compression is to provide bunch lengths that satisfy \( \sigma_z \ll \lambda/(2\pi) \) (where \( \lambda \) is the wavelength of the RF wave associated with the subsequent accelerating
structure). The latter condition insures that no significant LPS quadratic distortion is imparted during acceleration in the subsequent linac. For the L-band linac of ASTA this sets the upper requirement $\sigma_{\perp} \leq 800 \, \mu m$ \cite{4}. For a Gaussian distribution this requirement corresponds to a $I \approx 500 \, A$ at $Q = 3.2 \, nC$ and $C \sim 3.7 \, m^{-1}$ resulting in the tolerable bending-plane emittance dilution as shown in Tab.4, along with compression to $\sigma_{\perp} \approx 800 \, \mu m$ for each of the other bunch charges. Relative emittance growth is significantly smaller for 1-nC bunch charge and below.

6.2. High peak current production

Despite their relatively poor transverse emittance, fully compressed bunches could be used to generate copious amounts of radiation via a given electromagnetic process. The spectral-angular fluence emitted by a bunch of $N \gg 1$ electrons from any electromagnetic process is related to the single-electron spectral fluence, $\frac{d^2W}{d\Omega dt^2}$, via

$$\left.\frac{d^2W}{d\Omega dt^2}\right|_{\nu} \approx \left.\frac{d^2W}{d\Omega dt^2}\right|_{\nu} [N + N^2|S(\omega)|^2],$$  \hspace{1cm} (3)

where $\omega \equiv 2\pi f$ ($f$ is the frequency) and $S(\omega)$, the bunch form factor (BFF), is the intensity-normalized Fourier transform of the normalized charge distribution $S(t)$ \cite{29}. The former equation assumes the bunch can be approximated as a line charge distribution and is practically valid as long as the rms bunch duration $\sigma_{t}$ and transverse size $\sigma_{\perp}$ satisfy $\sigma_{\perp} \ll \sigma_{t}/\gamma$ where $\gamma$ is the Lorentz factor and $c$ is the velocity of light. When the BFF approaches unity, $\left.\frac{d^2W}{d\Omega dt^2}\right|_{\nu} \propto N^2$ and the radiation is termed “coherent radiation”.

At ASTA the availability of a superconducting linac coupled with a non-integer radiation-generation mechanism (e.g. diffraction radiation \cite{20}) could lead to the production of single-cycle THz pulses repeated at 3 MHz over 1-ms. As an example we consider the worst-case scenario of a fully compressed 3.2-nC bunch; the dependency of the BFF over frequency appears in Fig.15 (left plot). The BFF starts to take off at frequency lower than $f \approx 1 \, THz$ thereby supporting the generation of coherently-enhanced radiation at these frequencies. A limitation might come from the large transverse emittance that would prevent the beam to be focused to a transverse...
spot RMS size below the required \( \sigma_{\perp}/\gamma \) value. However a statistical analysis indicates that the central part of the beam containing approximately 15% of the beam population (or 500-pC out of the original 3.2-nC bunch) has emittances below 10 \( \mu \text{m} \) resulting in beam \( \sigma_{\perp} \leq 100 \mu \text{m} \); see Fig. 15 (right plot). It should be pointed out that the lower-charge cases investigated in the previous section would result in shorter pulses with associated BFFs that contain higher-frequency content (see also Fig. 5).

### 6.3. Compressed flat-beam generation

An important asset of the ASTA photoinjector is its capability to generate beams with high-transverse emittance ratios known as flat beams. Immersing the photocathode in a magnetic field introduces a canonical angular momentum \( \langle L \rangle = eB_{0}\sigma_{z} \), with \( B_{0} \) the magnetic field on the photocathode surface, and \( \sigma_{z} \) the RMS transverse size of the drive-laser spot on the photocathode [31]. As the beam exits the solenoidal field provided by lenses L1 and L2, the angular momentum is purely kinetic resulting in a beam coupled in the two transverse planes. Three skew quadrupoles in the beamline can apply the torque necessary to cancel the angular momentum [32, 33]. As a result, the final beam’s transverse emittance partition is given by

\[
\epsilon_{x,z} = \epsilon_{y} \left( \frac{\sigma_{y}}{2\beta\gamma L} \right),
\]

where \( \epsilon_{y} \) is the normalized uncorrelated emittance of the magnetized beam prior to the transformer, \( \beta \) and \( \gamma \) the Lorentz factors, \( L \equiv \langle L \rangle / 2p_{z} \), and \( p_{z} \) the longitudi-
nal momentum. Note that the product $e_{x,i}e_{y,i} = (e_{y})^2$. If compressed these flat beams may have applications in Smith-Purcell FELs [34] or for beam-driven acceleration techniques using asymmetric structures [35]. It may also be possible to mitigate the emittance growth in BC1 by having a beam that is wide in the direction of the chicane bend.

In this section, we explore the behavior of flat beams in the low-energy bunch compressor at ASTA, for the different initial emittance ratios $\rho \equiv \varepsilon_{x,i}/\varepsilon_{y,i}$. In order to produce these bunches, we took the 3.2-nC bunch presented earlier and numerically scaled the macroparticle coordinates to produce the desired transverse emittance ratios while constraining the product $e_{x,i}e_{y,i} = 5^2 \mu m^2$. Due to the large transverse aspect ratio of the bunches, the criterion given in Eq.1 is generally not satisfied and it is therefore anticipated that the projected CSR model is inadequate, thus we use CSRtrack’s P2P model to simulate the flat beams and neglect the 1DP model. The parameters used for flat beam simulations follow those used in the previous section, with the exception of the macroparticle horizontal size used in the CSRtrack P2P model. Due to the much greater transverse dimension we set $\sigma_h = 0.2$ mm. In addition, IMPACT-z SC+CSR simulations were also performed to evaluate the emittance growth in the vertical plane. The simulated emittance growth is shown in Fig.16 for a 3.2-nC bunch with an initial LPS chirp of $C = 5.2$ m$^{-1}$. As expected the relative emittance dilution is reduced as the initial emittance ratio $\rho$ increases. The agreement between CSRtrack and IMPACT-z for the bending-plane emittance dilution is remarkable (within $\sim 30\%$) given the large transverse horizontal beam sizes. IMPACT-z predicts that the vertical emittance increases by a factor of 1.5 to 1.8 over the range of considered initial emittance ratios $\rho \in [1, 500]$. The four-dimensional transverse emittance growth $\varepsilon_4 \equiv \sqrt{\varepsilon_x^2}$ is mitigated for the larger initial flat-beam emittance ratios.

6.4. Double-bunch generation

The production of shaped electron bunches has a large number of applications including the investigation of wakefield and beam-driven acceleration techniques. By operating the low energy bunch compressor with LPS chirp $C > 5.5$ m$^{-1}$ leads to over compression and results in a structured longitudinal charge distribution. Figure 17 confirms, for the case of $Q = 3.2$ nC, that a bi-modal distribution could be generated with a separation between its peaks ($\sim 300$ $\mu$m) consistent with requirements from beam-driven acceleration such as plasma-wakefield and dielectric-wakefield acceleration techniques. In addition the distance between the peaks could be controlled to some degree by slight changes over the initial LPS chirp. The full-bunch and slice-at-peak-current horizontal emittance at $C \sim 5.5$ m$^{-1}$ are 67.0 and 75.1 $\mu$m, respectively, compared to 106 and 107 $\mu$m for the maximum-compression case ($C \sim 5.2$ m$^{-1}$). These bending-plane normalized emittances of $\sim 75$ $\mu$m can still be focused to a sub-mm or sub-100-$\mu$m transverse spot size at respectively $\sim 40$ MeV and $\sim 250$ MeV (the latter energy corresponds to acceleration of the 40 MeV beam into one of ASTA accelerating cryomodules).
7. Summary

In this paper we presented numerical studies of a low-energy magnetic bunch compressor similar to the one being installed at the ASTA facility at Fermilab. Our results indicate that low-energy compression can be a viable path as a first stage compressor for a multistage compression scheme. In addition the capability of such a low energy compressor to provide high-peak currents at low energy (≈ 40 MeV) could have important applications when combined with the long-macropulse capability of the ASTA’s superconducting accelerating module such as, e.g. the production of single-cycle THz radiation from diffraction radiation. As part of our investigation we used several computer programs and observed an acceptable agreement when simulating the evolution of the LPS distributions during compression. The simulated transverse-phase-space parameters downstream of the bunch compressor have discrepancies that are inherent to the capabilities of each of the model but provide similar emittance growth and trade-off curves. Based on these observations, simple (and faster) models could be used to optimize the bunch compression design with a quick turn around while first-principle model could provide high-fidelity simulations of the optimized designs. The installation and commissioning of the BC1 bunch compressor along with the available diagnostics at ASTA will provide a unique experimental platform for benchmarking the simulation codes currently available in regimes where the CSR and SC effects can play similar roles.

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