WASP-30b: A 61 $M_{\text{Jup}}$ BROWN DWARF TRANSITING A $V = 12$, F8 STAR

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We report the discovery of a 61-Jupiter-mass brown dwarf (BD), which transits its F8V host star, WASP-30, every 4.16 days. From a range of age indicators we estimate the system age to be 1–2 Gyr. We derive a radius (0.89 ± 0.02 $R_{\text{Jup}}$) for the companion that is consistent with that predicted (0.914 $R_{\text{Jup}}$) by a model of a 1 Gyr old, non-irradiated BD with a dusty atmosphere. The location of WASP-30b in the minimum of the mass–radius relation is consistent with the quantitative prediction of Chabrier & Baraffe, thus confirming the theory.

Key words: binaries: eclipsing – brown dwarfs – stars: individual (WASP-30)

Online-only material: color figures, machine-readable table

1. INTRODUCTION

A brown dwarf (BD) is traditionally defined as an object with a mass above the deuterium-burning limit (13 $M_{\text{Jup}}$, e.g., Chabrier et al. 2000a) and below the hydrogen-burning limit (0.07 $M_{\odot}$; e.g., Chabrier et al. 2000b). However, an alternative suggestion is that the manner in which an object forms should determine whether it is a planet or a BD. Thus, if an object formed by core accretion of dust and ices in a protoplanetary disk then it would be a planet, and if it formed by gravoturbulent collapse of a molecular cloud, as do stars, then it would be a BD.

Studies such as the Caballero et al. (2007) observations of a young open cluster core find a continuous mass function down to ~ 6 $M_{\text{Jup}}$, indicating that the star formation mechanism can produce objects with planetary masses. This is supported by theoretical studies (Padoan & Nordlund 2004; Hennebelle & Chabrier 2008) which suggest that gravothermal fragmentation of molecular clouds produces stars and BDs down to a few Jupiter masses in numbers comparable to the observationally determined distribution. In contrast, when taking into account planetary migration through the protoplanetary disk, the core accretion process might result in giant planets with masses of up to 10 $M_{\text{Jup}}$ (Alibert et al. 2005) or even 25 $M_{\text{Jup}}$ (Mordasini et al. 2008). Sahlmann et al. (2010) see evidence for a bimodal distribution in BD masses, with the less-massive group presumably representing the high-mass tail of the planetary distribution.

An accurate, precise measurement of an object’s radius is therefore required to probe for the existence of a core and thus discriminate between the two formation mechanisms. For example, the radius of the 8 $M_{\text{Jup}}$ body, HAT-P-2b, is consistent with an irradiated planet incorporating a 340-Earth-mass core, but is smaller than if it were coreless (Leconte et al. 2009). The 22 $M_{\text{Jup}}$ CoRoT-3b (Deleuil et al. 2008) is sufficiently massive to qualify as a BD under the traditional definition, but the radius of this object is uncertain at the 7% level. This is higher than the 3% required to discriminate between the absence or the presence of a core and thus determine how it formed (Leconte et al. 2009).

We report the discovery of a 61-Jupiter-mass brown dwarf (BD), which transits its F8V host star, WASP-30, every 4.16 days. From a range of age indicators we estimate the system age to be 1–2 Gyr. We derive a radius (0.89 ± 0.02 $R_{\text{Jup}}$) for the companion that is consistent with that predicted (0.914 $R_{\text{Jup}}$) by a model of a 1 Gyr old, non-irradiated BD with a dusty atmosphere. The location of WASP-30b in the minimum of the mass–radius relation is consistent with the quantitative prediction of Chabrier & Baraffe, thus confirming the theory.

There is less ambiguity around the upper end of the BD mass regime: if a body is sufficiently massive to fuse hydrogen then it is a star, otherwise it is a BD. High-mass BDs with precise radius measurements are useful for testing BD evolution models, as it is in the high-mass regime that models predict the greatest changes in radius with age (e.g., Baraffe et al. 2003). Stassun et al. (2006) discovered a BD eclipsing binary system in the Orion Nebula star-forming region, with masses of 57 ± 5 $M_{\text{Jup}}$ and 36 ± 3 $M_{\text{Jup}}$. With very large radii of 0.699 ± 0.034 $R_{\odot}$ and 0.511 ± 0.026 $R_{\odot}$, it seems that these objects are in the earliest stages of gravitational contraction. Similar to the NLTT 41135 system, LHS 6343 C (Johnson et al. 2010; J. A. Johnson 2010, private communication) is a 63 $M_{\text{Jup}}$ BD that transits one member of an M-dwarf binary system. In this case, the transits are full and the radius (0.825 ± 0.023 $R_{\text{Jup}}$) of this object is precisely determined. CoRoT-15b (Bouchy et al. 2010) is a 63 $M_{\text{Jup}}$ mass BD in a 3 day orbit around an F7V star. Due to the faintness of the host star ($V$ ~ 16), the BD radius (1.12$^{+0.30}_{-0.15}$ $R_{\text{Jup}}$) is not yet well determined.

To test and refine models of BD formation and evolution, a population of well-characterized objects is required. In this Letter, we present the discovery of WASP-30b, a 61 $M_{\text{Jup}}$ BD that transits its moderately bright host star.

2. OBSERVATIONS

WASP-30 is a $V = 11.9$, F8V star located in Aquarius, on the border with Cetus. A transit search (Collier Cameron et al. 2006) of WASP-South data from 2008 July to November found a strong, 4.16 day periodicity. Further observations in 2009 with
both WASP instruments (Pollacco et al. 2006) led to a total of 17,612 usable photometric measurements (Figure 1).

Using the CORALIE spectrograph mounted on the 1.2 m Euler-Swiss telescope (Baranne et al. 1996; Queloz et al. 2000), we obtained 16 spectra of WASP-30 in 2009. Radial velocity (RV) measurements were computed by weighted cross-correlation (Baranne et al. 1996; Pepe et al. 2005) with a numerical G2-spectral template. RV variations were detected with the same period found from the WASP photometry and with a semi-amplitude of 6.6 km s\(^{-1}\), consistent with a substellar-mass companion. The RV measurements are listed in Table 1 and are plotted in Figure 1.

To test the hypothesis that the RV variations are due to spectral line distortions caused by a blended eclipsing binary or starspots, we performed a line-bisector analysis (Queloz et al. 2001) of the CORALIE cross-correlation functions. The lack of correlation between bisector span and RV (Figure 1) supports our conclusion that the periodic dimming of WASP-30 and its RV variations are due to the sub-stellar orbiting body, WASP-30b.

To refine the system parameters, we obtained high signal-to-noise transit photometry through a Gunn \(r\) filter with the Euler-Swiss telescope on 2010 August 1 (Figure 1; Table 2). The data were affected by light cloud and a guiding issue caused by the close proximity of the bright Moon (69% illumination, 26° from target).

### Table 1

Radial Velocity Measurements of WASP-30

| BJD−2,450,000 | RV (km s\(^{-1}\)) | \(\sigma_{\text{RV}}\) (km s\(^{-1}\)) | BS (km s\(^{-1}\)) |
|---------------|-----------------|-----------------|-----------------|
| 5009.9065     | 14.275          | 0.032           | 0.077           |
| 5040.8722     | 1.298           | 0.050           | −0.190          |
| 5092.6977     | 14.348          | 0.044           | 0.048           |
| 5095.6894     | 5.163           | 0.030           | −0.039          |
| 5096.5476     | 12.979          | 0.040           | −0.020          |
| 5096.8712     | 14.452          | 0.041           | 0.005           |
| 5097.5351     | 12.509          | 0.041           | 0.052           |
| 5097.8735     | 9.732           | 0.048           | 0.039           |
| 5098.5538     | 3.350           | 0.040           | −0.006          |
| 5113.5209     | 14.451          | 0.031           | −0.127          |
| 5113.5877     | 14.563          | 0.036           | −0.008          |
| 5113.6111     | 14.515          | 0.038           | −0.009          |
| 5113.6343     | 14.582          | 0.034           | −0.009          |
| 5113.6576     | 14.488          | 0.031           | −0.032          |
| 5113.6809     | 14.335          | 0.030           | −0.009          |
| 5113.7041     | 14.508          | 0.031           | −0.061          |

### Table 2

Euler Photometry of WASP-30

| BJD−2,450,000 | Relative Flux | \(\sigma_{\text{flux}}\) |
|---------------|---------------|-----------------|
| 5409.693659   | 1.00123       | 0.00185         |
| 5409.695116   | 0.99944       | 0.00185         |
| 5409.696482   | 1.00213       | 0.00185         |
| ...           | ...           | ...             |
| 5409.917840   | 1.00020       | 0.00185         |
| 5409.918650   | 0.99791       | 0.00185         |

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

3. STELLAR PARAMETERS

The 16 CORALIE spectra of WASP-30 were co-added to produce a single spectrum with a typical signal-to-noise ratio of 70:1. The analysis was performed using the methods given in Gillon et al. (2009). The H\(_\alpha\) line was used to determine the effective temperature (\(T_{\text{eff}}\)), while the Na \(I\) D and Mg \(\text{i}\)
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Table 3
System Parameters for WASP-30

| Parameter                                      | Value       |
|------------------------------------------------|-------------|
| Stellar Parameters from Spectroscopic Analysis |             |
| R.A. = 23 h 53 m 38.03 s, decl. = −10° 07′ 05″.1 (J2000) |             |
| TYC 5834-95-1, 2MASS 23533805-1007049             |             |
| Teff (K)                                        | 6100 ± 100  |
| log g                                          | 4.3 ± 0.1   |
| ξi (km s⁻¹)                                    | 1.1 ± 0.1   |
| v sin i (km s⁻¹)                               | 14.2 ± 1.1  |
| [Fe/Fe]                                        | 0.08 ± 0.10 |
| [Si/Fe]                                        | 0.04 ± 0.13 |
| [Ca/Fe]                                        | 0.10 ± 0.14 |
| [Ti/Fe]                                        | −0.01 ± 0.14|
| [Ni/Fe]                                        | −0.08 ± 0.13|
| A11                                            | 2.87 ± 0.10 |

Parameters from MCMC Analysis

| Parameter                                      | Value       |
|------------------------------------------------|-------------|
| P (days)                                       | 4.156736 ± 0.000013 |
| Tc (HJD)                                       | 245533.94879 ± 0.000076 |
| T14 (days)                                     | 0.1595 ± 0.00017 |
| T12 = T14 (days)                               | 0.01060 ± 0.000225 |
| ΔF = Rp²/R∗²                                  | 0.00498 ± 0.00017 |
| b                                              | 0.066 ± 0.024 |
| i (°)                                          | 89.57 ± 0.28 |
| K1 (km s⁻¹)                                    | 6.627 ± 0.015 |
| a (AU)                                         | 0.05325 ± 0.00039 |
| e                                               | 0 (adopted) |
| y' (km s⁻¹)                                    | 7.9177 ± 0.0099 |
| Me (M⊙)                                        | 1.166 ± 0.026 |
| Re (R⊙)                                        | 1.295 ± 0.019 |
| log g, (cgs)                                   | 4.280 ± 0.010 |
| log ρ, (ρ)                                    | 0.537 ± 0.019 |
| Teff (K)                                       | 6201 ± 97 |
| [Fe/Fe]                                        | −0.03 ± 0.10 |
| Mr (M_Jup)                                     | 60.96 ± 0.89 |
| Rp (R_Jup)                                     | 0.889 ± 0.021 |
| log g, (cgs)                                   | 5.247 ± 0.019 |
| ρp (ρ)                                        | 86.8 ± 5.7 |
| T_eq (K)                                       | 1474 ± 25 |

b lines were used as surface gravity (log g,) diagnostics. The parameters obtained from the analysis are given in the top panel of Table 3. The elemental abundances were determined from equivalent width measurements of several clean and unblended lines. A value for microturbulence (ξi) was determined from Fe i using the method of Magain (1984). The quoted error estimates include that given by the uncertainties in ξi, log g, and ξi, as well as the scatter due to measurement and atomic data uncertainties. Our quoted lithium abundance takes account of non-local thermodynamic equilibrium corrections (Carlsson et al. 1994), with a value of A11 = 2.95 resulting when neglecting them.

The projected stellar rotation velocity (v sin i) was determined by fitting the profiles of several unblended Fe i lines. A value for macroturbulence (v_max) of 4.7 ± 0.3 km s⁻¹ was assumed, based on the tabulation by Gray (2008), and an instrumental FWHM of 0.11 ± 0.01 Å, determined from the telluric lines around 6300 Å. A best-fitting value of v sin i = 14.2 ± 1.1 km s⁻¹ was obtained.

4. SYSTEM PARAMETERS

The WASP and Euler photometry were combined with the CORALIE RV measurements in a simultaneous Markov Chain Monte Carlo (MCMC) analysis (Collier Cameron et al. 2007; Pollacco et al. 2008). Our proposal parameters are: Teff, P, ΔF, T14, b, K1, Teff, [Fe/H], v sin i, and cos i, sin i (Collier Cameron et al. 2007; Enoch et al. 2010). Here, Teff is the epoch of mid-transit, P is the orbital period, ΔF = Rp²/R∗² is the fractional flux-deficit that would be observed during transit in the absence of limb-darkening, T14 is the total transit duration (from first to fourth contact), b is the impact parameter of the BD’s path across the stellar disk, K1 is the semi-amplitude of the stellar reflex velocity, Teff is the stellar effective temperature, [Fe/H] is the stellar metallicity, e is the orbital eccentricity, and ω is the argument of periastron.

As Ford (2006) notes, it is convenient to use e cos ω and e sin ω as MCMC jump parameters, because these two quantities are nearly orthogonal and their joint probability density function is well-behaved when the eccentricity is small and ω is highly uncertain. Ford cautions, that the use of e cos ω and e sin ω as jump parameters implicitly imposes a prior on the eccentricity that increases linearly with e. Instead we use √e cos ω and √e sin ω as jump parameters, which restores a uniform prior on e.

At each step in the MCMC procedure, each proposal parameter is perturbed from its previous value by a small, random amount. From the proposal parameters, model light and RV curves are generated and χ² is calculated from their comparison with the data. A step is accepted if χ² (our merit function) is lower than for the previous step, and a step with higher χ² is accepted with probability exp(−Δχ²/2). In this way, the parameter space around the optimum solution is thoroughly explored.

The value and uncertainty for each parameter are respectively taken as the median and central 68.3% confidence interval of the parameter’s marginalized posterior probability distribution.

From the proposal parameters, we calculate the mass M, radius R, density ρ, and surface gravity log g of the star (which we denote with the subscript *) and the planet (which we denote with subscript P). At each step, the stellar density is measured from the transit light curve (Seager & Mallén-Ornelas 2003). This is input in to the empirical mass calibration of Torres et al. (2010), as modified by Enoch et al. (2010), to obtain an estimate of the stellar mass. We also calculate the equilibrium temperature of the planet T_eq, assuming it to be a blackbody with efficient redistribution of energy from the planet’s day side to its night side, the transit ingress and egress durations, T12 and T14, and the orbital semi-major axis a.

With eccentricity floating, we find e = 0.0021 ± 0.0024. Applying the “F-test” of Lucy & Sweeney (1971), we find a 70% probability that the fitted eccentricity could have arisen by chance if the underlying orbit is in fact circular. As such, we impose a circular orbit, but we note that doing so has no significant effect on the solution in this case.

Without exquisite photometry, our implementation of MCMC tends to bias the impact parameter, and thus Rp and Rp, to higher values. This is because, with low signal-to-noise photometry, the transit ingress and egress durations are uncertain, and symmetric uncertainties in those translate into asymmetric uncertainties in b and thus in Rp. We therefore place a main-sequence (MS) prior on the star, which is reasonable given the star’s apparent age (Section 5). With the MS prior, a Bayesian penalty ensures that, in accepted MCMC steps, the values of stellar radius are consistent with the values of stellar mass for a MS star (Collier Cameron et al. 2007).

The median parameter values and their 1σ uncertainties from our MCMC analysis are presented in the bottom panel of Table 3. The corresponding transit light curves and RV curve are shown.
in Figure 1. When not imposing an MS prior, the best-fitting values we obtain are $b = 0.24^{+0.24}_{-0.16}$, $R_\text{u} = 1.33^{+0.14}_{-0.02} R_\odot$, and $R_P = 0.925^{+0.118}_{-0.040} R_{\text{Jup}}$.

5. SYSTEM AGE AND COMPANION RADIUS

The high lithium abundance ($A_{\text{Li}} = 2.87 \pm 0.10$) found in WASP-30 implies an age most likely between that of open clusters such as $\alpha$ Per (50 Myr; $A_{\text{Li}} = 2.97 \pm 0.13$) and the Hyades (600 Myr; $A_{\text{Li}} = 2.77 \pm 0.21$), and almost certainly younger than 2 Gyr old open clusters such as NGC 752 ($A_{\text{Li}} = 2.65 \pm 0.13$; Sestito & Randich 2005).

Assuming aligned stellar-spin and planetary-orbit axes, the measured $v \sin i$ of WASP-30 and its derived stellar radius indicate a rotational period of $P_{\text{rot}} = 4.6 \pm 0.4$ days. After removing the transits from the WASP light curves, we searched them for evidence of rotational modulation. Though there are periodogram peaks at periods of 4.1 days, 4.3 days, and 4.7 days, the signal amplitudes are small. We thus conclude that there is no evidence of rotational modulation in the WASP-30 light curves, commensurate with expectations based on the star’s spectral type. The 4.6 day stellar rotation period is very close to the companion’s orbital period, suggesting that the two may be synchronized, thus preventing a gyrochronological age determination (Barnes 2007). The synchronization timescale (Zahn 1977) for the star is $0.23 \pm 0.02$ Gyr.

We searched within 15' of the sky position of WASP-30 for common proper motion stars. The $V = 13.6$ star USNO-B1.0 0800-0674908 is 13/09 away and appears to be comoving with WASP-30. Using 2MASS photometry, we constructed a color–magnitude diagram. A distance modulus of $8.50 \pm 0.05$ (500 $\pm$ 10 pc) was required to place the comoving star on the MS and suggests that WASP-30 is $\sim$2.3 Gyr old (Figure 2, upper panel). However, the apparent magnitude and spectral type of WASP-30 suggest a smaller distance modulus of $7.9 \pm 0.2$ (366 $\pm$ 77 pc). As the comoving star may be a mere line-of-sight neighbor, this age determination should be treated with caution.

In the lower panel of Figure 2, WASP-30b is plotted in a mass–radius diagram together with isochrones of isolated BDs from models with dusty atmospheres (DUSTY00; Chabrier et al. 2000b) and models with dust-free atmospheres (COND03; Baraffe et al. 2003). By an age of 1 Gyr, an isolated BD with the mass of WASP-30b is expected to have cooled to $T_{\text{eq}} \sim 1700$ K, and the transition from dusty L-dwarfs to dust-free T-dwarfs is expected to take place at $T_{\text{eq}} = 1300$–1700 K (Baraffe et al. 2003). Considering this and the fact that WASP-30b is highly irradiated, it is likely that the dusty atmosphere models of Chabrier et al. (2000b) are more representative. Depending on the transition temperature, WASP-30b may remain a dusty L-dwarf for the lifetime of its host star or it may at some point transition to a dust-free T-dwarf.

WASP-30b has a mass of $0.05819 \pm 0.00084 M_\odot$ and a radius of $0.0914 \pm 0.0022 R_\odot$. The DUSTY00 models predict the radius of an isolated, 0.06 $M_\odot$ BD to be 0.149, 0.104, 0.094, and 0.084 $R_\odot$ at ages of, respectively, 0.1, 0.5, 1, and 5 Gyr. From a simple linear interpolation of the DUSTY00 model values, the measured radius of WASP-30b and its 1$\sigma$ uncertainty suggests its age is 2 $\pm$ 1 Gyr. The measured radius is inconsistent with the DUSTY00 model for a 0.5 Gyr BD at the 5.7$\sigma$ level and inconsistent with a 5 Gyr BD at the 3.4$\sigma$ level, but consistent with a 1 Gyr BD at the 1.2$\sigma$ level. This agreement would be slightly better if the DUSTY00 models took account of irradiation. Baraffe et al. (2003) find that the irradiation of a dust-free atmosphere, at the level of irradiation experienced by WASP-30b, results in radii larger by 10% for a 1 $M_\text{Jup}$ planet and larger by 7% for a 10 $M_\text{Jup}$ planet. However, they note that the evolution of dust-free atmospheres are more affected by irradiation than are dusty atmospheres, and WASP-30b is considerably more massive.

WASP-30’s lithium abundance favors a young system age of 50–600 Myr, though lithium is a poor age indicator for an F8 star, and an age of up to 2 Gyr is not ruled out. The apparent rotational synchronization of the host star places a lower limit of 50–600 Myr, though the system may have been born synchronized) and a possible companion star suggests an older age of 2.3 Gyr. Given the BD’s measured radius, the DUSTY00 BD model indicates an age of 2 $\pm$ 1 Gyr. Taken together, we suggest that an age of 1–2 Gyr is most likely.

6. DISCUSSION

The discovery of WASP-30b heralds the first unambiguous observational determination of the mass–radius relation (MRR) in the BD regime, and so we have added BDs to white dwarfs...
and neutron stars in the list of quantum-dominated objects with radius determinations.

Chabrier & Baraffe (2000) performed the first quantitative theoretical calculation of the MRR in the substellar and low-mass-star domain, predicting a minimum in the MRR at high BD masses (see Section 3.1 of that paper). The location of WASP-30b in the MRR minimum is consistent with the quantitative prediction of Chabrier & Baraffe (2000), thus confirming the theory.

Thus far, we know of two other high-mass BDs that transit stars: CoRoT-15b (63 $M_{\odot}$; Bouchy et al. 2010) and LHS 6343 C (63 $M_{\odot}$; Johnson et al. 2010; J. A. Johnson 2010, private communication). The radius of CoRoT-15b (1.12$^{+0.30}_{-0.15}$ $R_{\odot}$) is uncertain and the age of the system is currently unconstrained. The faintness of the host star ($V \sim 16$) makes improving this situation difficult. LHS 6343 C was found to transit one member of an M-dwarf binary system using Kepler photometry (KIC 10002261; e.g., Borucki et al. 2010). It initially seemed that LHS 6343 C was larger than predicted for a BD of its mass and age (Johnson et al. 2010). However, after a re-evaluation of the treatment of the third light in the system, it seems to be consistent (J. A. Johnson 2010, private communication).

Chabrier et al. (2000b) predict that BDs with $M < 0.05 M_\odot$ do not burn lithium, those with $M > 0.06 M_\odot$ burn essentially all their lithium by an age of 0.5 Gyr, and those with an intermediate mass ($M = 0.055 M_\odot$) burn half their lithium by 0.5 Gyr, two-thirds by 1 Gyr, and three-quarters by 5 Gyr. With a mass of 0.0582 ± 0.0008 and an age of 1–2 Gyr, WASP-30b is likely to have burned most of its supply of lithium.

WASP-30b has the second smallest companion-to-star size ratio ($\Delta F = R_{\text{companion}}^2/R_{\star}^2 = 0.0050$) of all sub-stellar bodies so far discovered by ground-based transit surveys. The star in the system with the smaller size ratio, HAT-P-11 ($V = 9.6; \Delta F = 0.0033$; Bakos et al. 2010), is eight times brighter than WASP-30. As it is far easier to find such an object around a smaller, cooler star, the discovery of WASP-30b suggests that high-mass, sub-stellar objects in short orbits around cooler stars are rare.

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