Total kinetic energy release in the fast neutron induced fission of $^{232}\text{Th}$ and $^{235}\text{U}$

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We have measured the total kinetic energy release (TKE), its variance and associated fission product distributions for the neutron induced fission of $^{232}\text{Th}$ and $^{235}\text{U}$ for $E_n = 2$–90 MeV. The neutron energies were determined on an event by event basis by time of flight measurements with the white spectrum neutron beam from LANSCE. The TKE decreases non-linearly with increasing neutron energy for both systems, while the TKE variances are sensitive indicators of $n^{th}$ chance fission. The associated fission product distributions show the decrease in TKE with increasing beam energy that is due to the increasing probability of symmetric fission— which has a lower associated TKE—and the decreasing TKE associated with asymmetric fission, presumably due to the decreasing importance of the $A \sim 132$ shell structures.

Keywords: Fission; TKE; variance of TKE; fission product mass distributions; $^{232}\text{Th}(n,f)$; $^{235}\text{U}(n,f)$

1. Introduction

We report the total kinetic energy release (TKE), the variance ($\sigma^2$) of the TKE distribution, and the associated mass yield distributions in the fast neutron induced fission of $^{232}\text{Th}$ and $^{235}\text{U}(E_n = 2$–90 MeV). (A preliminary version of the $^{235}\text{U}$ study has been reported before.) While these data have implications for nuclear weapons and reactors, our primary goal is to better understand the fission process. As we increase the energy input to these systems from $\sim 0$ to $\sim 90$ MeV, we find that the primary energy released in fission, the TKE, decreases by a few MeV (6.0 and 7.4 MeV, for $^{232}\text{Th}$ and $^{235}\text{U}$, respectively). Clearly the dissipation of the incident neutron energy in the fissioning system takes place on a time scale ($\sim 10^{-21}$ s) that is short compared to the times for the large scale collective motion of fission ($\sim 10^{-18}$ s). The dissipated energy appears in the excitation energy of the fission fragments while the small decrease in the observed TKE reflects changes in the fission process in these reactions. Understanding
and characterizing these processes will be the central theme of this work. By studying the $^{232}$Th(n,f) and the $^{235}$U(n,f) processes simultaneously, we hope to compare and contrast the physics of the different fissioning systems upon the fission product mass and energy distributions.

2. Experimental

This experiment was carried out at the Weapons Neutron Research Facility (WNR) at the Los Alamos Neutron Science Center (LANSCE) at the Los Alamos National Laboratory $^{2,3}$ over a 5 day period in December, 2014. “White spectrum” neutron beams were generated from an unmoderated tungsten spallation source using the 800 MeV proton beam from the LANSCE linear accelerator. The measured white spectrum near the target position is shown in Fig. 1. The proton beam is pulsed allowing one to measure the time of flight (energy) of the neutrons arriving at the experimental area. Note the beam intensities of $\sim 10^5$ n/s—a number that is higher than the beam used at radioactive beam facilities—but substantially less than that normally used in fission studies. Thus one has to run for long periods of time and/or to use large solid angle detection systems.

![Fig. 1. Measured neutron beam intensity near the target location on the 15R beam line during this experiment](image)

The detailed description of this experiment has been reported previously $^1$. We refer the reader to $^1$ for details. In $^1$, assuming $\nu_{\text{tot}} = 2.43$ and
\[ \nu_{pre} = 0.0^4 \], we found that the measured thermal neutron induced post-fission TKE was \( E_{TKE^{th,n}} = 169.8 \pm 0.4 \text{ MeV} \) and that \( \sigma_{TKE^{th}} = 103.6 \pm 0.2 \text{ MeV} \). The pre-fission TKE is \( E_{TKE^{*,th}} = 170.7 \pm 0.4 \text{ MeV} \). This result is consistent with the previous measurements of \( 171.9 \pm 1.4^5 \), \( 172.0 \pm 2.0^6 \) and \( 170.1 \pm 0.5^7 \). For the spontaneous fission of \(^{252}\text{Cf}\), we found the average post fission light fragment TKE was \( 103.2 \pm 1.1 \text{ MeV} \) while the average heavy fragment TKE was \( 78.6 \pm 1.5 \text{ MeV} \). This is to be compared to post neutron emission values\(^8\) for the light fragment of \( 103.8 \pm 0.5 \text{ MeV} \) and heavy fragment energy of \( 79.4 \pm 0.4 \text{ MeV} \). **Our measurements of the TKE should be regarded as absolute measurements.**

3. Results and Discussion

The TKE data for each system was binned into neutron energy bins. The resulting TKE distributions for each neutron energy bin were found to be Gaussian in shape with no observed deviations from a Gaussian shape. The mean values of the TKE for each bin for each system are shown in Figure 2. The mean TKE values decrease non-linearly with increasing neutron energy in a manner that can be represented by a polynomial in neutron energy. These data are in good agreement with previous measurements at lower energies\(^9-13\). The Viola systematics\(^14\), which applies only to pre-neutron emission data, would predict TKE values of \( 170.1 \) and \( 163.8 \text{ MeV} \), respectively, for the \(^{235}\text{U}(n,f)\) and \(^{232}\text{Th}(n,f)\) reactions. That “Viola ratio” \((170.1/163.8)\) is remarkably similar to the observed TKE ratio for the two systems, perhaps indicating the same underlying physics. This observation of a non-linear decrease in TKE with increasing excitation energy is certainly not unique—as it was observed in \(^{235}\text{U}(p,f)\) reactions decades ago by Ferguson et al.\(^15\). Ferguson et al. attributed this decrease to an increase in symmetric fission (which has a lower TKE) and the possible “washing out” of the \( A=132 \) shell structure with increasing beam energy although they could not resolve the two contributing effects. Using a modal analysis, Brosa et al.\(^16\) attempted to resolve this issue for these systems \((^{232}\text{Th}(n,f)\) and \(^{235}\text{U}(n,f))\) for neutron energies of \( E_n = 0 - 6 \text{ MeV} \). They concluded that the reduced TKE as a function of \( E_n \) could be attributed simply to a decrease in symmetric fission. A similar result using an “empirical fission potential” was achieved by by Wang et al.\(^17\). The choice of this low neutron energy regime is not accidental in that it assures that only single chance fission is being studied.

The variances of the TKE distributions for the \(^{235}\text{U}\) reaction are shown
in Figure 3 along with the unpublished data of Duke. It is striking to note the occurrence of peaks in the variances at the same neutron energies in both works and the correspondence of those peaks with the onset of nth chance fission. The variances of the TKE distributions represent a sensitive tool for detecting the occurrence of multiple chance fission.

The pre-neutron emission fission product mass yield distributions for the two reactions are shown in Figures 4 and 5. Uncertainties in the individual points can be deduced from the Y(A) values which are the number of counts in a given mass bin. For both systems, one observes increasing amounts of symmetric fission with increasing neutron energy. The 235U distributions contain about 75K events while the 232Th distributions contain about 20K events so the uncertainties are smaller in the 235U data, but because of the shapes of the distributions, the symmetric component is more noticeable in the 232Th data.

In Figure 6(a), we show the evolution with neutron energy of the mass sorted TKE values for the 232Th(n,f) reaction. One sees the TKE associated with symmetric fission is roughly constant with increasing E_n with a lower average TKE value. The TKE associated with asymmetric fission decreases with increasing E_n which is incompatible with the notion that overall TKE
decrease with increasing $E_n$ is due \textbf{solely} to the increased probability of symmetric fission. Further confirmation of this trend is seen in Figure 6 (b) where one can observe the steep changes in the TKE associated with the $A \sim 134$ fragments (associated with the spherical shell structures at $Z=50, N=82$, which accompany the trend of the symmetric fission TKE with increasing $E_n$.

To learn about fission, one would like to compare these measured properties of the TKE distributions with predictions of nuclear theory. The theories should give “predictive” estimates of measured quantities rather than “post-dictive” estimates. Two semi-empirical models\textsuperscript{18,19} for the TKE release as a function of neutron energy are compared to our data in Figure 7. Both models take into account multiple chance fission. The predictions of the GEF code differ from the data above $E_n = 20$ MeV while the predictions of\textsuperscript{19} agree quite well with the data. One should note that more “fundamental” calculations of fission product yields\textsuperscript{20} do exist for these systems. In these calculations, there is a washing out of the shell effect and changes in the pairing energies with increasing excitation energy. The
Fig. 4. The measured pre-neutron emission fission product mass distributions as a function of neutron energy for the $^{235}\text{U}(n,f)$ reaction.

calculation$^{20}$ is limited to predicting (correctly) the fission product mass distribution for $^{235}\text{U}(n,f)$ at $E_n=14.5$ MeV.

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Fig. 5. (a) The measured pre-neutron emission fission product mass distributions as a function of neutron energy for the $^{232}$Th(n,f) reaction.

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