INTEGER HOMOLOGY 3-SPHERES ADMIT IRREDUCIBLE REPRESENTATIONS IN $SL(2, \mathbb{C})$

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Abstract. We prove that the fundamental group of any integer homology 3-sphere different from the 3-sphere admits irreducible representations of its fundamental group in $SL(2, \mathbb{C})$. For hyperbolic integer homology spheres this comes with the definition, and for Seifert fibered integer homology spheres this is well known. We prove that the splicing of any two non-trivial knots in $S^3$ admits an irreducible $SU(2)$-representation. By work of Boileau, Rubinstein, and Wang, the general case follows.

Using a result of Kuperberg, we get the corollary that the problem of 3-sphere recognition is in the complexity class $\text{coNP}$, provided the generalised Riemann hypothesis holds.

To prove our result, we establish a topological fact about the image of the $SU(2)$-representation variety of a non-trivial knot complement into the representation variety of its boundary torus, a pillowcase. For this, we use holonomy perturbations of the Chern-Simons function in an exhaustive way – we show that any area-preserving self-map of the pillowcase fixing the four singular points, and which is isotopic to the identity, can be $C^0$-approximated by maps which are realised geometrically through holonomy perturbations of the flatness equation in a thickened torus. To conclude, we use a stretching argument in instanton gauge theory, and a non-vanishing result of Kronheimer and Mrowka for Donaldson’s invariants of a 4-manifold which contains the 0-surgery of a knot as a splitting hypersurface.

Introduction

A lot of the facts known about 3-manifolds come from knowledge about their fundamental groups and the representations of these groups in fairly simple Lie groups like $SO(3)$, $SU(2)$ and $SL(2, \mathbb{C})$. Particular interest has
been paid to integer homology spheres, as these turn out to be of particular interest, having trivial abelianisation. For instance, the $SU(2)$-representation varieties of Seifert-fibred homology spheres have been described quite explicitly by Fintushel and Stern in terms of linkages \cite{21}, using the explicit descriptions of the fundamental groups of such 3-manifolds. Hyperbolic integer homology 3-spheres come with irreducible representations of their fundamental group in $SL(2,\mathbb{C})$, because the projectivisation $PSL(2,\mathbb{C}) = SL(2,\mathbb{C})/\{\pm \text{id}\}$ is the orientation-preserving isometry group of hyperbolic 3-space.

In general, it is quite difficult to describe these representation varieties. A prominent result in this direction is Kronheimer and Mrowka’s proof of the Property $P$ conjecture of knots in $S^3$. They have established \cite{31, 32} that for a fraction $p/q$ such that $|p/q| \leq 2$ the $p/q$-surgery on a non-trivial knot in $S^3$ always admits irreducible representations of its fundamental group in $SU(2)$. Their proof is indirect, using highly non-trivial instanton gauge theory. More recently, Baldwin and Sivek have related the existence of irreducible $SU(2)$-representations to Stein fillings, see \cite{5, 6}.

One of the main results that we establish is the following

**Theorem 9.4.** Let $Y$ be an integer homology 3-sphere different from the 3-sphere. Then there is an irreducible representation $\rho: \pi_1(Y) \to SL(2,\mathbb{C})$.

Hence this is a characterisation of the 3-sphere among integer homology 3-spheres.

Problem 3.105 in Kirby’s problem list asks whether any integer homology 3-sphere different from the 3-sphere admits an irreducible representation in $SU(2)$. Theorem 9.4 answers this question for the group $SL(2,\mathbb{C})$.

Using results of Boileau, Rubinstein and Wang on the theory of dominations of 3-manifolds by degree-1 maps, see Theorem \ref{thm:domination} below, Theorem 9.4 follows from the following theorem that we establish.

**Theorem 8.3.** Let $K$ and $K'$ be two non-trivial knots in $S^3$. Then there is an irreducible representation $\rho: \pi_1(Y_{K,K'}) \to SU(2)$ from the fundamental group of the splicing $Y_{K,K'}$ of $K$ and $K'$.

Here the splicing is defined as follows: Given two knots $K, K'$ in $S^3$ with tubular neighbourhoods $N(K)$ and $N(K')$, the splicing is the integer homology 3-sphere

$$Y_{K,K'} = Y(K) \cup_{\phi} Y(K'),$$
obtained by glueing the knot complements $Y(K) = S^3 \setminus N(K)\complement$ and $Y(K') = S^3 \setminus N(K')\complement$ along their boundary tori by a map $\varphi$ which maps a meridian of $K$ to a longitude of $K'$, and vice versa.

Quite a few facts are known about the representation variety
\[ R(K) = \text{Hom}(\pi_1(Y(K)), SU(2))/SU(2). \] (1)
of the knot complement and its image in the representation variety of the boundary torus. For instance, for non-trivial knots $K$ there are always irreducible representations, and these have to contain 1-dimensional semi-algebraic varieties. However, it is not clear a priori whether the two images from $R(K)$ and $R(K')$ intersect in the representation variety of the splitting torus of $Y_{K,K'}$, and hence yield an irreducible representation of $\pi_1(Y_{K,K'})$. (It is well-known, for instance, that this holds if $K$ and $K'$ are non-trivial torus knots.)

Our proof of Theorem 8.3 will rely on a topological result (Theorem 7.1) about the image of $R(K)$ in the representation variety of the boundary torus that we will establish for any non-trivial knots. This topological result ensures that the images of the two representation varieties $R(K)$ and $R(K')$ in the representation variety of the boundary torus intersect in a representation of the boundary torus and hence give rise to an irreducible representation of $\pi_1(Y_{K,K'})$.

Before stating this result, we make a few general comments about Theorem 9.4.

**Remark.** A statement analogous to Theorem 9.4 does not hold for the more general class of rational homology 3-spheres, see [38].

**Remark.** Theorem 9.4 would hold with the group $SU(2)$ at the place of $SL(2, \mathbb{C})$ if one could show that any integer homology 3-sphere which is hyperbolic admits irreducible $SU(2)$-representations of its fundamental group.

It isn’t so clear whether one can expect that this holds. For instance, a corresponding statement fails for some hyperbolic rational homology 3-spheres:

**Remark.** By Cornwell’s results in [12], the branched double cover $\Sigma_2(8_{18})$ of the knot $8_{18}$ does not admit irreducible $SU(2)$-representations. The rational homology sphere $\Sigma_2(8_{18})$ is hyperbolic. In fact, the knot $8_{18}$ is the Turk’s head knot with notation $(2 \times 4)^* $ in [3] Section 4.3, and in this reference it is shown that its branched double cover is hyperbolic.
Theorem 9.4 is a result about 3-manifolds, and not about groups. In fact, there are superperfect groups which admit no non-trivial representations in $SL(2, \mathbb{C})$, see Proposition 10.2. By a result of Kervaire which uses higher-dimensional surgery theory, all of these appear as fundamental groups of homology spheres of dimension $n \geq 5$. Those which admit a presentation of deficiency 0 even appear as the fundamental group of a 4-dimensional homology sphere. Hence we obtain the following result, contrasting Theorem 9.4.

**Theorem 10.1.** In any dimension $n \geq 4$ there is an integer homology sphere $W^n$ different from the $n$-sphere $S^n$ such that the fundamental group $\pi_1(W)$ has only the trivial representation in $SL(2, \mathbb{C})$.

Theorem 9.4 also has the following application on the complexity of the problem of 3-sphere recognition, due to work of Kuperberg [34]. (We are thankful to John Baldwin and Steven Sivek for turning our attention to this reference.)

**Theorem 11.3.** Let $Y$ be an integer homology 3-sphere, described by a Heegaard diagram. Then the assertion that $Y$ is not the 3-sphere lies in the complexity class $NP$, provided the generalized Riemann hypothesis holds.

The pillowcase and the $SU(2)$-representation variety of knot complements. The space of $SU(2)$-representations of the fundamental group of a two-dimensional torus $T^2$ modulo conjugation,

$$R(T^2) = \text{Hom}(\mathbb{Z}^2, SU(2))/SU(2)$$

is homeomorphic to the pillowcase, a 2-dimensional sphere. In fact, if we denote generators of $\pi_1(T^2) \cong \mathbb{Z}^2$ by $m$ and $l$, then for a representation $\rho$ we may suppose that

$$\rho(m) = \begin{bmatrix} e^{i\alpha} & 0 \\ 0 & e^{-i\alpha} \end{bmatrix}, \quad \text{and} \quad \rho(l) = \begin{bmatrix} e^{i\beta} & 0 \\ 0 & e^{-i\beta} \end{bmatrix},$$

and hence we can associate to $\rho$ a pair $(\alpha, \beta) \in [0, 2\pi] \times [0, 2\pi]$, which we also can think of as being a point on the two-dimensional torus $T = \mathbb{R}^2/2\pi\mathbb{Z}^2$. However, it is easily seen that a representation to which we associate $(2\pi - \alpha, 2\pi - \beta)$ is conjugate to $\rho$. This is the only ambiguity, however, as the trace of an element in $SU(2)$ determines its conjugacy class. Therefore $R(T^2)$ is isomorphic to the quotient of the torus $T$ by the
hyperlattice involution $\tau: (\alpha, \beta) \mapsto (-\alpha, -\beta)$. This has four fixed points, and its quotient

$$R(T^2) = T/\tau$$

is homeomorphic to a two-dimensional sphere. It can also be seen as the quotient of the fundamental domain $[0, \pi] \times [0, 2\pi]$ for $\tau$ by identifications on the boundary as indicated in Figure 1 below.

\begin{figure}[h]
\centering
\includegraphics[width=0.5\textwidth]{pillowcase}
\caption{The glueing pattern for obtaining the pillowcase from a rectangle, and the image of the representation variety $R(K)$ of the trefoil in the pillowcase.}
\end{figure}

Having a non-trivial knot $K$ in $S^3$, the 3-manifold $Y(K) = S^3 \setminus \mathcal{N}(K)$ obtained by removing a tubular neighbourhood of $K$ from $S^3$ is a 3-manifold with boundary a two-dimensional torus. Associated is the representation variety $R(K)$, defined above.

We can restrict representations from $R(K)$ to its boundary torus, thereby obtaining a map $R(K) \to R(T^2)$ into the pillowcase. Figure 1 above shows the image of $R(K)$ when $K$ is the trefoil, once in the pillowcase, and once in the fundamental domain $[0, \pi] \times [0, 2\pi]$. Here the first coordinate corresponds to $\rho(m_K)$, where $m_K$ is a meridian to the knot $K$, and the second coordinate corresponds to $\rho(l_K)$, where $l_K$ is a longitude of the knot $K$.

All abelian representations map to the red line ‘at the bottom’. In fact, as $l_K$ is a product of commutators in the fundamental group of the knot complement, an abelian representation has to map $l_K$ to the identity, and hence its image in $R(T^2)$ lies on the bottom line $\{\beta = 0 \mod 2\pi \mathbb{Z}\}$, and for any $\alpha$ we can find an abelian representation corresponding to $(\alpha, 0)$. 
As a meridian normally generates the knot group, the two lines \( \{ \alpha = 0 \mod 2\pi \mathbb{Z} \} \) and \( \{ \alpha = \pi \mod 2\pi \mathbb{Z} \} \) only contain the two central representations from \( R(K) \), both with \( \{ \beta = 0 \mod 2\pi \mathbb{Z} \} \).

If we cut the pillowcase open along these two lines we obtain a cylinder \( C = [0, \pi] \times \mathbb{R}/2\pi \mathbb{Z} \) (in the glueing pattern of Figure 1 this means that we do not perform the identifications along the four indicated vertical boundary lines.) There has been some evidence that the image of \( R(K) \) in the cylinder \( C \) always contains a closed curve which is homologically non-trivial in the cylinder. For instance, it is always the case for non-trivial torus knots. Furthermore, the author has seen images of \( R(K) \) in the pillowcase that have been determined numerically by Culler [13], where this property was also satisfied in all examples. One of our main result states that this is always the case.

![Figure 2](image.png)

**Figure 2.** Can there be an embedded path \( c \) from \( P \) to \( Q \) which does not intersect the image of the representation variety of the knot (in red), as in the picture?

**Theorem 7.1.** Let \( K \) be a non-trivial knot in \( S^3 \). Then the image of \( R(K) \) in the cut-open pillowcase \( C = [0, \pi] \times (\mathbb{R}/2\pi \mathbb{Z}) \) contains an embedded curve which is homologically non-trivial in \( H_1(C; \mathbb{Z}) \cong \mathbb{Z} \).

Let us observe that if this result holds, then there cannot be a curve like the curve \( c \) in Figure 2 connecting the points \( P \) and \( Q \), and which does not intersect the image of \( R(K) \).

Conversely, if we can show that any embedded path from \( P \) to \( Q \) as in Figure 2 disjoint from the line \( \{ \beta = 0 \mod 2\pi \mathbb{Z} \} \) (where the reducibles
of any \( R(K) \) map, does intersect \( R(K) \), then Theorem 7.1 follows by an argument using Alexander duality, and the fact that the image of \( R(K) \) in \( R(T^2) \) is a compact embedded graph (see Lemma 7.3).

**Theorem 7.2.** Let \( K \) be a non-trivial knot. Then any embedded path from \( P = (0, \pi) \) to \( Q = (\pi, \pi) \) in the pillowcase, missing the line \( \{ \beta = 0 \mod 2\pi \mathbb{Z} \} \), has an intersection point with the image of \( R(K) \).

**Remark 7.3.** To the best of the author’s knowledge, there has been no previous constraint on the image of \( R(K) \) in the pillowcase that would contradict a picture like in Figure 2. The immersions in the shape of an ‘8’ take into account results by Herald [25] about the area that can be enclosed by the image of a circle or irreducibles in \( R(K) \).

**Geometric realisation of isotopies by holonomy perturbations.** To establish Theorem 7.2 we use holonomy perturbations of the flatness equation of a Hermitian rank-2 bundle on the 0-surgery on the knot \( K \) in an exhaustive way. At the heart of the argument is a technical result which is of independent interest. We state it here in a simplified way which is sufficient to understand the strategy of the proof of the main results.

**Theorem 4.2.** Let \( (\psi_t)_{t \in [0, 1]} : T \to T \) be an isotopy of the torus \( T = \mathbb{R}^2/2\pi \mathbb{Z}^2 \) (which we think of as the branched double cover of the pillowcase), such that \( \psi_0 = \text{id} \), and let \( \varepsilon > 0 \) be given. Let \( M = [0, 1] \times S^1 \times S^1 \) be a thickened torus, and let \( E \to M \) be the trivial Hermitian rank-2 bundle. Then there is a map \( \theta^t \), defined on the space of all \( SU(2) \)-connections on \( E \), with values in \( \Omega^2(M; \text{su}(E)) \), (the space of 2-forms with values in the adjoint bundle,) with compact support in the interior of \( M \), such that any connection \( A \) of the perturbed flatness equation

\[ F_A = \theta^t(A) \]

is reducible, and such that the holonomies at the two boundaries are related in the following way:

\[ \text{Hol}(A)|_{(1) \times S^1 \times S^1} = \phi_t(\text{Hol}(A)|_{(0) \times S^1 \times S^1}) \]

where \( \phi_t \) is \( \varepsilon \)-close to \( \psi_t \). Here we consider \( \text{Hol}(A)|_{i \times S^1 \times S^1} \) for \( i = 0, 1 \) as a point in the pillowcase determined by the holonomies of the two curves \( m_0 = \{0\} \times S^1 \times pt \) and \( l_0 = \{0\} \times pt \times S^1 \), and likewise for \( m_1 \) and \( l_1 \).

An essential inspiration to a strategy of proof of Theorem 4.2 is due to a discussion that the author had with Frank Kutzschebauch, from whom he
learned about the role of shearing maps in a domain of complex geometry now known as Andersen-Lempert theory [2, 3], see also [27].

Holonomy perturbations of the flatness equation have been introduced by Floer [22], and since then used by various authors, mainly in an approach to transversality, see for instance work of Taubes, Donaldson or Herald [45, 16, 25], or, more recently, Herald-Kirk [26]. Big perturbations as in our situation have been considered by Braam-Donaldson [10] for establishing the surgery exact triangle in instanton Floer homology, by Kronheimer-Mrowka in [32], and more recently by Lin in [35, 36].

**Strategy of proof of Theorem 7.2.** Suppose there were an embedded path from \( P \) to \( Q \), missing \( R(K) \). We lift the problem to the branched double cover of the pillowcase, the torus \( T = \mathbb{R}^2/2\pi\mathbb{Z}^2 \). We can find a closed embedded curve \( c \) which is invariant under the hyperelliptic involution, starting at the fixed point \( \hat{P} \), and passing through the fixed point \( \hat{Q} \), and which misses the double cover of the image of \( R(K) \) in \( T \). There is an isotopy \( \varphi_t \) from \( \varphi_0 = \text{id} \) such that curves \( c = \varphi_1(c_0) \), where \( c_0 \) is the straight line \( \{ \beta = \pi \} \) from \( \hat{P} \) to \( \hat{P} \) via \( \hat{Q} \), and we may assume that the curves \( c_t = \varphi_t(c_0) \) do not pass through the line \( \{ \beta = 0 \} \), where the reducibles map. By a lemma that we learned from Thomas Vogel, we can realise the isotopy \( \varphi_t \) by an isotopy through area-preserving maps \( \psi_t : T \to T \), with \( \psi_0 = \text{id} \).

Elements of \( R(K) \) that map to the line \( c_0 \) are those \( SU(2) \)-representations of the fundamental group of the knot complement that extend to \( SO(3) \)-representations of the fundamental group of the 0-surgery \( Y_0(K) \) which cannot be lifted to \( SU(2) \)-representations. We understand these as flat \( SO(3) \)-connections on an \( SO(3) \)-bundle over \( Y_0(K) \), or equivalently, as critical points of the Chern-Simons function \( CS \). Our main technical result Theorem 4.2 allows us to define a holonomy perturbation \( \Phi(t) \) of the Chern-Simons function, such that the critical points of \( CS + \Phi(t) \) correspond to points of \( R(K) \) which lie on the line \( \phi_t(c_0) \), where \( \phi_t \) can be chosen arbitrarily close to \( \psi_t \), and hence \( \phi_t(c_0) \) arbitrarily close to the curve \( c_t \).

Our assumption means that the Chern-Simons function \( CS + \Phi(1) \) has no critical points. But this yields a contradiction, if we use a non-vanishing theorem of Donaldson’s invariants of a symplectic 4-manifold \( X \) containing the 0-surgery \( Y_0(K) \) as a separating hypersurface, due to Kronheimer and Mrowka. In fact, we perform a neck-stretching argument to show that if \( CS + \Phi(1) \) has no critical point, then Donaldson’s invariant of \( X \) must vanish, following the strategy of Kronheimer-Mrowka in [32]. We deal with big holonomy perturbations here, and so we have to deal with the critical points
of the family CS + Φ(t), as well as with a corresponding one-parameter family of holonomy perturbations on X with some care. (For instance, the usual argument to avoid reducibles in the instanton moduli spaces with only perturbations of the metric fails in our situation.)

**Strategy of Proof of Theorem 4.2.** For the most simple holonomy perturbations of the flatness equation $F_A = 0$ one can choose in $[0, 1] \times S^1 \times S^1$, restrictions of the holonomy to the two boundary components are related by a shearing map between the two pillowcases appearing as the representation varieties of the two boundaries. One can also realise iterations of shearing maps by a finite nested family of such holonomy perturbations. In order to prove Theorem 4.2 one therefore has to show that the group generated by shearing diffeomorphisms is $C^0$-dense in the group of area-preserving diffeomorphisms of the pillowcase. Actually, we need a slightly more refined version: We have to show that any isotopy through area-preserving diffeomorphisms can be $C^0$-approximated by an isotopy through shearings.

From an area-preserving isotopy we pass to a time-dependent divergence-free vector field. The isotopy is then given by the ‘flow’ of this vector field. We approximate the flow of the time-dependent vector field by a composition of flows of time-independent vector fields. For these we choose the time-dependent vector field at consecutive times, differing by just a small amount of time. We then take the Fourier series corresponding to these time-dependent vector fields. A Fourier term turns out to be a vector field of shearing type – its flow is an isotopy through shearing maps. So we first approximate the consecutive time-independent vector fields by finite Fourier series, and we estimate the difference of the corresponding flows. Finally, we approximate the flow of a finite Fourier sum by flows along the single Fourier terms.

**Outline of the paper.** Section 1 contains a review of the Chern-Simons function and Floer’s holonomy perturbation of it. Section 2 describes nested families of holonomy perturbations on a thickened torus and the maps obtained between the pillowcases of the two boundaries. Section 3 contains the crucial approximation results. Section 4 just contains the main technical result which puts the results of Section 2 and 3 together. Section 5 discusses the holonomy perturbations on the 0-surgery of a knot, and it describes the perturbed $SU(2)$-representation variety in terms of representations of the knot complement with an imposed boundary condition. Section 6 deals with Donaldson’s invariants, Kronheimer-Mrowka’s non-vanishing
result, and with instanton moduli spaces on 4-manifolds which contain holonomy perturbations in a neck. Section 7 proves Theorems 7.1. Section 8 proves Theorem 8.3, our result about splicing of two non-trivial knots. Section 9 deals with general integer homology 3-spheres. Section 10 deals with the higher dimensional case, and Section 11 deals with the complexity of 3-sphere recognition.

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1. A class of holonomy perturbations of the Chern-Simons function

Let $Y$ be a closed Riemannian 3-manifold, and let $E \to Y$ be a Hermitian bundle of rank 2 with determinant line bundle $w \to Y$. Let $\theta$ be a connection in the line bundle $w$. We denote by $\mathcal{A}$ the affine space of Hermitian connections in $E$ which induce the connection $\theta$ in the determinant line bundle. This space of connections is the same as the space of $SO(3)$-connections in the associated bundle $\mathfrak{su}(E)$. The group of automorphisms of $E$ with determinant 1 is called the gauge group and is denoted by $\mathcal{G}$. It acts on $\mathcal{A}$ in a natural way. A connection is called reducible if its stabiliser under the $\mathcal{G}$-action is different from $\mathbb{Z}/2 \cong \pm \text{id}$, and irreducible otherwise.

We are interested in the critical points of the Chern-Simons function

$$\text{CS}: \mathcal{A} \to \mathbb{R}$$

$$A \mapsto \int_Y \text{tr}((A - A_0) \wedge ((F_A)_0 - (F_{A_0})_0)),$$

where $A_0$ is some fixed reference connection in $\mathcal{A}$, $F_A$ denotes the curvature of the connection $A$, and $(F_A)_0$ denotes the trace-free part of the curvature. This trace-free part of the curvature is equal to the curvature of the connection, when thought of being the corresponding $SO(3)$-connection in the associated bundle $\mathfrak{su}(E)$, and tr denotes the trace on $2 \times 2$ matrices. This
function is invariant under the action of the gauge group.

The critical points of the Chern-Simons function CS correspond to connections $A$ satisfying the equation

$$(F_A)_0 = 0.$$ 

**Definition 1.1.** We denote by

$$R^w(Y) = \{ [A] \in \mathcal{A}/\mathcal{G} | (F_A)_0 = 0 \} \quad (2)$$

the space of equivalence classes of critical points. Via the holonomy, this space corresponds to conjugacy classes of representations $\rho: \pi_1(Y) \to SO(3)$ with second Stiefel-Whitney class $w_2(\rho) \equiv w \mod 2$, see [32, Lemma 4]. Therefore, $R^w(Y)$ is also called the $SO(3)$-representation variety of $Y$ associated to $w$.

For what we have in mind, we will need a perturbed version of the Chern-Simons function, and hence for the flatness equations. The perturbation we will use have been introduced by Floer [22] and are called holonomy perturbations.

Let $\chi: SU(2) \to \mathbb{R}$ be a smooth function which only depends on the conjugacy classes of $SU(2)$, acted on by itself via the conjugation action. Any element in $SU(2)$ is conjugate to a diagonal element, and hence there is a $2\pi$-periodic even function $g: \mathbb{R} \to \mathbb{R}$ such that

$$\chi\left(\begin{bmatrix} e^{it} & 0 \\ 0 & e^{-it} \end{bmatrix}\right) = g(t)$$

for all $t \in \mathbb{R}$. Let furthermore $\Sigma$ be a compact surface with boundary, and let $\mu$ be a real-valued 2-form which has compact support in the interior of $\Sigma$ and with $\int_{\Sigma} \mu = 1$. Let $\iota: \Sigma \times S^1 \to Y$ be an embedding. Let $N \subseteq Y$ be a codimension-0 submanifold containing the image of $\iota$, and such that the bundle $E$ is trivialised over $N$ in such a way that the connection $\theta$ in $\text{det}(E)$ induces the trivial product connection in the determinant line bundle of our trivialisation of $E$ over $N$. This means that connections in $\mathcal{A}$ can be understood as $SU(2)$-connections in $E$ restricted to $N$.

Associated to this data, we can define a function

$$\Phi: \mathcal{A} \to \mathbb{R}$$

which is invariant under the action of the gauge group $\mathcal{G}$. For $z \in \Sigma$, we denote by $\iota_z: S^1 \to Y$ the circle $t \mapsto \iota(z,t)$. A connection $A \in \mathcal{A}$
being an SU(2)-connection over the image of \( \iota \), the holonomy \( \text{Hol}_{\iota_z}(A) \) of \( A \) around the loop \( \iota_z \) (with variable starting point) is a section of the bundle of automorphism of \( E \) with determinant 1 over the loop. Since \( \chi \) is a class function, \( \chi(\text{Hol}_{\iota_z}(A)) \) is well-defined. We can therefore define
\[
\Phi(A) = \int_{\Sigma} \chi(\text{Hol}_{\iota_z}(A)) \mu(z). \quad (3)
\]

Later in this paper, we will consider a finite sequence of such embeddings, all supported in a submanifold \( M \) of codimension 0. For some \( n \in \mathbb{N} \), let \( \iota_k: S^1 \times \Sigma \to M \subseteq Y \) be a sequence of embeddings for \( k = 0, \ldots, n - 1 \) such that the image of \( \iota_k \) is disjoint from the image of \( \iota_l \) for \( k \neq l \). (The surface \( \Sigma_k \) could also be taken to be dependent on \( k \), but we will not need this in our situation.)

We also suppose class functions \( \chi_k: SU(2) \to \mathbb{R} \) corresponding to even, 2\( \pi \)-periodic functions \( g_k: \mathbb{R} \to \mathbb{R} \) as above to be chosen, for \( k = 0, \ldots, n - 1 \), and we continue to assume that \( \mu \) is a 2-form on \( \Sigma \) with support in the interior of \( \Sigma \) and integral 1. Just as what lead to (3), we obtain a finite sequence of functions
\[
\Phi_k: \mathcal{A} \to \mathbb{R}, \quad k = 0, \ldots, n - 1.
\]

The following is essentially proved in [10].

**Proposition 1.2.** The critical points of the perturbed Chern-Simons function
\[
\text{CS} + \sum_{k=0}^{n-1} \Phi_k: \mathcal{A} \to \mathbb{R}
\]
are the elements \( A \in \mathcal{A} \) which solve the equation
\[
(F_A)_0 = \sum_{k=0}^{n-1} \chi'_k(\text{Hol}_{\iota_k}(A)) \mu_k, \quad (4)
\]
where \( \chi'_k \) is an equivariant map \( SU(2) \to \mathfrak{su}(2) \) which is the dual to the derivative of \( \chi_k \) with respect to the Killing form on \( \mathfrak{su}(2) \), and where the pull-back of \( \mu_k \) by \( \iota_k \) is the 2-form on \( S^1 \times \Sigma \) which is obtained by pulling back \( \mu \) from \( \Sigma \) to \( S^1 \times \Sigma \).

**Definition 1.3.** For choices made as above, we denote
\[
R^w_{\{\iota_k, \chi_k\}}(Y) = \{ [A] \in \mathcal{A}/\mathcal{G} \mid (F_A)_0 = \sum_{k=0}^{n-1} \chi'_k(\text{Hol}_{\iota_k}(A)) \mu_k \},
\]
and we call this the perturbed $SO(3)$-representation variety of $Y$ associated to $w$ and the holonomy perturbation data $\{\iota_k, \chi_k\}$.

**Remark 1.4.** This definition depends also on the choice of a trivialisation of the bundle $E$ over some codimension-0 submanifold which contains the images of the $\iota_k$.

## 2. Nested holonomy perturbations on a thickened torus

In this subsection we place ourselves in the situation of a particular 3-manifold, namely $M = [0, 1] \times S^1 \times S^1$, a thickened torus. We suppose $E \to M$ is the trivial $SU(2)$-bundle. We will now study a family of holonomy perturbations with nested support in $M$. Such kind of perturbations have also been studied recently by Herald and Kirk [26], but in a different perspective.

Let $\Sigma = [0, 1] \times S^1$ be the 2-dimensional annulus. For $k = 0, \ldots, n - 1$, let

$$A_k = \begin{pmatrix} a_k & c_k \\ b_k & d_k \end{pmatrix} \in SL(2, \mathbb{Z}),$$

be a sequence of matrices. We consider the sequence of embeddings $\iota_k: \Sigma \times S^1 \to M$ given by

$$\iota_k: [0, 1] \times S^1 \times S^1 \to M$$

$$\left( t, \begin{pmatrix} z \\ w \end{pmatrix} \right) \mapsto \begin{pmatrix} k + t \\ n \end{pmatrix}, \begin{pmatrix} a_k & c_k \\ b_k & d_k \end{pmatrix} \begin{pmatrix} z \\ w \end{pmatrix}, \quad (5)$$

where we understand $S^1 = \mathbb{R}/\mathbb{Z}$, so that the matrix multiplication is understood in the usual sense.

For $k = 0, \ldots, n - 1$, we also suppose class functions $\chi_k: SU(2) \to \mathbb{R}$ corresponding to even, $2\pi$-periodic functions $g_k: \mathbb{R} \to \mathbb{R}$ determined by

$$\chi_k \left( \begin{bmatrix} e^{it} & 0 \\ 0 & e^{-it} \end{bmatrix} \right) = g_k(t), \quad (6)$$

and we assume that $\mu$ is a 2-form on $\Sigma$ with support in the interior of $\Sigma$ and integral 1. We denote by $f_k = g'_k$ the derivative of $g_k$, so that $f_k$ is a $2\pi$-periodic odd function.
Let $z_0 = (x_0, y_0) \in S^1 \times S^1$ be some chosen base point. For $k = 0, \ldots, n$ we will consider the closed curves

$$m_k = \left\{ \frac{k}{n} \right\} \times S^1 \times \{ y_0 \} \quad \text{and} \quad l_k = \left\{ \frac{k}{n} \right\} \times \{ x_0 \} \times S^1. \quad (7)$$

**Proposition 2.1.** Let $A$ be an SU(2)-connection on the trivial bundle over the thickened torus $[0, 1] \times S^1 \times S^1$ satisfying the perturbed flatness equation

$$F_A = \sum_{k=0}^{n-1} \chi_k' \left( \text{Hol}_{l_k}(A) \right) \mu_k, \quad (8)$$

associated to the holonomy perturbation data $\{ l_k, \chi_k \}$ determined by equations (5) and (6) above. Then $A$ is reducible, and up to a gauge-transformation we may suppose that $A$ respects the splitting $E = M \times (\mathbb{C} \oplus \mathbb{C})$. If we write the holonomies along the curves $m_k, l_k$ as

$$\text{Hol}_{m_k}(A) = \begin{bmatrix} e^{i\alpha_k} & 0 \\ 0 & e^{-i\alpha_k} \end{bmatrix}, \quad \text{and} \quad \text{Hol}_{l_k}(A) = \begin{bmatrix} e^{i\beta_k} & 0 \\ 0 & e^{-i\beta_k} \end{bmatrix}, \quad (10)$$

for $k = 0, \ldots, n$, then we have the relationships

$$\begin{pmatrix} \alpha_{k+1} \\ \beta_{k+1} \end{pmatrix} = \begin{pmatrix} \alpha_k \\ \beta_k \end{pmatrix} + f_k(-b_k \alpha_k + a_k \beta_k) \begin{pmatrix} a_k \\ b_k \end{pmatrix} \quad (9)$$

for $k = 0, \ldots, n - 1$. Here $f_k$ is the $2\pi$-periodic odd function which is the derivative of the function $g_k$ associated to the class function $\chi_k$ by equation (6) above.

**Proof.** The connection $A$ satisfying equation (8) must be reducible. This is proved in [10, Lemma 4]. Up to a gauge transformation, we may therefore suppose that

$$\text{Hol}_{m_k}(A) = \begin{bmatrix} e^{i\alpha_k} & 0 \\ 0 & e^{-i\alpha_k} \end{bmatrix},$$

$$\text{Hol}_{l_k}(A) = \begin{bmatrix} e^{i\beta_k} & 0 \\ 0 & e^{-i\beta_k} \end{bmatrix}, \quad (10)$$

in the trivial bundle $E$ over $M$, for $k = 0, \ldots, n$. The pairs $(\alpha_k, \beta_k) = (\alpha_k(A), \beta_k(A))$ are determined by $A$ up to sign, and up to addition of some pair $(n, m) \in 2\pi\mathbb{Z} \times 2\pi\mathbb{Z}$, as other such choices correspond to equal or conjugate holonomies.
The annulus Σ has oriented boundary given by the two embedded circles
\( m_\pm : S^1 \to \partial \Sigma \), where \( m_\pm \) runs around \( \{0\} \times S^1 \), and \( m_\pm \) runs around \( \{1\} \times S^1 \). Similarly, we denote by \( l_\pm \) the curve \( \{1\} \times \{x_0\} \times S^1 \), and by \( l_\pm \) the curve \( \{0\} \times \{x_0\} \times S^1 \). We suppose that \( m_\pm \) corresponds to the boundary orientation, and that \( m_\pm \) corresponds to its opposite. For the connection \( \bar{A} = \iota^*_k A \) we denote by
\[
\text{Hol}_{m_\pm}(\bar{A}) = \left[ e^{i\alpha_\pm} \ 0 \ 0 \ e^{-i\alpha_\pm} \right].
\]
the holonomy of the curves \( m_\pm \). For any curve \( l = \{p_0\} \times S^1 \) we denote by
\[
\text{Hol}_l(\bar{A}) = \left[ e^{i\beta} \ 0 \ 0 \ e^{-i\beta} \right],
\]
its holonomy. In fact, as \( A \) satisfies equation (8), this doesn’t depend on the choice of the point \( p_0 \in \Sigma \) by construction, see loc. cit.

By [10, Lemma 4], we must have
\[
\alpha_+ - \alpha_- = f_k(\beta),
\]
If we denote \( \beta_\pm := \beta \), we can write this relationship also in the form
\[
\chi_{\iota_k} : \left( \begin{array}{c} \alpha_- \\ \beta_- \end{array} \right) \mapsto \left( \begin{array}{c} \alpha_+ \\ \beta_+ \end{array} \right) = \left( \begin{array}{c} \alpha_- + f_k(\beta_-) \\ \beta_- \end{array} \right).
\] (11)

As the connection \( A \) is flat on the tori \( \{k\} \times S^1 \times S^1 \), its holonomy around loops in these tori only depends on the homology class of the loops. In particular, the holonomy of a loop homologous to \( pm_\pm + ql_\pm \) is given by
\[
\text{Hol}_{m_\pm}(A)^p \circ \text{Hol}_{l_\pm}(A)^q.
\]
Combining the map [11] together with the relationship of the curves on the tori \( \{\frac{k}{n}\} \times S^1 \times S^1 \), \( \{\frac{k+1}{n}\} \times S^1 \times S^1 \) with the curves \( m_\pm \) and \( l_\pm \) induced by \( \iota_k \), we obtain
\[
\left( \begin{array}{c} \alpha_{k+1} \\ \beta_{k+1} \end{array} \right) = (A_k \circ \chi_{\iota_k} \circ A_k^{-1}) \left( \begin{array}{c} \alpha_k \\ \beta_k \end{array} \right)
\] (12)
A straightforward computation now yields
\[
\left( \begin{array}{c} \alpha_{k+1} \\ \beta_{k+1} \end{array} \right) = \left( \begin{array}{c} \alpha_k \\ \beta_k \end{array} \right) + f_k(-b_k \alpha_k + a_k \beta_k) \left( \begin{array}{c} a_k \\ b_k \end{array} \right).
\] (13)
\[\square\]
The maps $\chi_{f_k}$ defined in (11) and the associated maps
\[ \zeta_k : = A_k \circ \chi_{f_k} \circ A_k^{-1} \] (14)
appearing in equation (12) are maps known as shearings.

**Definition and Proposition 2.2** (Shearing maps of the torus). Let $f : \mathbb{R} \rightarrow \mathbb{R}$ be a $2\pi$-periodic smooth function, and let $v \in \mathbb{R}^2$ be given. Let
\[ l : \mathbb{R}^2 \rightarrow \mathbb{R} \]
be a linear form which maps the lattice $2\pi \mathbb{Z}^2$ to $2\pi \mathbb{Z}$ and which contains the vector $v$ in its kernel. Then the map
\[ \phi : \mathbb{R}^2 \rightarrow \mathbb{R}^2 \]
\[ (\alpha, \beta) \mapsto (\alpha, \beta) + f(l((\alpha, \beta))) \cdot v \] (15)
descends to a well-defined diffeomorphism $T^2 \rightarrow T^2$ that we also denote by $\phi$, and that we call a shearing in direction $v \in \mathbb{R}^2$ associated to the function $f$ and the linear function $l$.

**Example 2.3.** The map $\zeta_k$ defined by equation (14) above is a shearing in direction $v = (a_k, b_k) \in \mathbb{Z}^2$ associated to the $2\pi$-periodic function $f_k$ and the linear form $l$ given by taking the standard inner product with the vector $w_k = (-b_k, a_k) \in \mathbb{Z}^2$ which is orthogonal to $v$,
\[ l : u \mapsto \langle u, w \rangle . \]

**Remark 2.4.** We notice that shearing maps are non-local, and therefore come with a certain rigidity: If a point $(\alpha, \beta) \in T^2$ is mapped by a shearing in direction $v \in \mathbb{R}^2$ to $(\alpha, \beta) + c \cdot v$ then any point $(\alpha, \beta) + d \cdot v$ is mapped to $(\alpha, \beta) + (d+c) \cdot v$. In other words, any point on the circle or line $(\alpha, \beta) + \mathbb{R} \cdot v$ is moved by the same amount $c$ in direction $v$ along that line. For instance, an isotopy with support in a small ball in $T^2$ can never be realised by a shearing.

Proposition 2.1 can be restated by saying that solutions $A$ to equation (8) have holonomies at the boundaries which are related by the finite sequence of shearings $\zeta_k$ defined in (14), up to a gauge transformation. There is a straightforward converse of Proposition 2.1.

**Proposition 2.5.** Given shearing maps $\phi_k : T^2 \rightarrow T^2$ in directions
\[ v_k = \begin{pmatrix} a_k \\ b_k \end{pmatrix} \in \mathbb{Z}^2 \]
and associated to 2π-periodic odd functions $f_k: \mathbb{R} \to \mathbb{R}$ and linear forms $l_k(u) = \langle u, w_k \rangle$ with $w_k = (-b_k, a_k)$ for $k = 0, \ldots, n-1$, there is some holonomy perturbation data $\{\iota_k, \chi_k\}_{k=0}^{n-1}$ such that the class functions $\chi_k$ associated to even 2π-periodic functions $g_k$ according to equation (6) satisfy

$$g'_k = f_k,$$

and such that the embeddings $\iota_k$ are related to matrices $A_k$ as in (5) above, with the following significance: If $[A]$ solves the perturbed flatness equation (8), then the resulting shearing maps $\zeta_k$, relating the holonomies as in equation (14), coincide with $\phi_k$, for $k = 0, \ldots, n-1$.

□

In the next section we will show that compositions of shearing maps are $C^0$-dense in the space of area-preserving maps of the 2-dimensional torus.

3. APPROXIMATION OF AREA-PRESERVING ISOTOPIES BY SHEARING ISOTOPIES

On the two-dimensional torus $T^2 = \mathbb{R}^2/\mathbb{Z}^2$ we have the hyperelliptic involution given by $\tau(x, y) = (-x, -y)$. It has four fixed points, and its quotient is the pillowcase, a topological 2-sphere with four distinguished points, the fixed points of the involution. When we speak of the $\mathbb{Z}/2$-action on $T^2$, we shall mean this involution.

**Definition 3.1.** Let $X$ be a vector field on $T^2$ that we naturally identify with a map $X: T^2 \to \mathbb{R}^2$. We shall say that $X$ is a vector field of shearing type if there is a direction $v \in \mathbb{R}^2$, a linear map $l: \mathbb{R}^2 \to \mathbb{R}$ which takes the lattice $2\pi \mathbb{Z}^2$ to $2\pi \mathbb{Z}$, and a smooth 2π-periodic function $f: \mathbb{R} \to \mathbb{R}$ such that we have

$$X(x, y) = f(l((x, y))) \cdot v$$

for all $(x, y)$.

**Remark 3.2.** If $X$ is a shearing vector field, then the associated map $(x, y) \mapsto (x, y) + X(x, y)$ is a shearing map, and conversely, every shearing map is of such type.

**Theorem 3.3.** Let

$$\psi: [0, 1] \times T^2 \to T^2$$

$$(t, (x, y)) \mapsto \psi_t(x, y)$$
be a $\mathbb{Z}/2$-equivariant smooth isotopy through area-preserving maps (which necessarily fixes the four fixed points of the hyperelliptic involution.) Then for any $\varepsilon > 0$, there is a $\mathbb{Z}/2$-equivariant map

$$\phi: [0, 1] \times T^2 \to T^2$$

$$(t, (x, y)) \mapsto \phi_t(x, y)$$

which is continuous in $t$, and smooth in $(x, y)$, such that we have

(i) $d(\psi_t(x, y), \phi_t(x, y)) < \varepsilon$

for all $t \in [0, 1]$ and all $(x, y) \in T^2$. Here $d$ denotes the metric on $T^2$ coming from the natural Euclidean structure.

(ii) For each $t \in [0, 1]$, the map $\phi_t$ is a finite composition of $\mathbb{Z}/2$-equivariant shearing maps, and there is a sequence $t_0, \ldots, t_{n+1} \in [0, 1]$ with $0 = t_0 < t_1 < \cdots < t_n < t_{n+1} = 1$ such that for all $i = 0, \ldots, n$ there is a $\mathbb{Z}/2$-equivariant shearing vector field $W_i: T^2 \to \mathbb{R}^2$ with direction $v \in \mathbb{Z}^2$, and we have

$$\phi_t(x, y) = \phi_{t_i}(x, y) + (t - t_i) \cdot W_i(x, y)$$

for $t_i \leq t \leq t_{i+1}$, and for all $(x, y) \in T^2$.

In other words, the smooth isotopy $(\psi_t)$ can be $C^0$-approximated by isotopies $(\phi_t)$ which are (in the coordinate $t$) piecewise isotopies through shearing isotopies.

The remainder of this section is devoted to the Proof of this theorem. The essential ingredient is the following Lemma.

**Lemma 3.4.** Let $X: T^2 \to \mathbb{R}^2$ be a divergence-free vector field,

$$\text{div}(X) = \frac{\partial X^1}{\partial x} + \frac{\partial X^2}{\partial y} \equiv 0.$$ 

Then for any $\varepsilon > 0$ there is a finite sequence of shearing vector fields $(X_i)_{i=0, \ldots, m}$ with directions $v_i \in \mathbb{Z}^2$ such that we have

$$\|X - \sum_{i=0}^m X_i\|_\infty < \varepsilon.$$

If $X$ is $\mathbb{Z}/2$-equivariant, then the vector fields $X_i$ can be chosen to be $\mathbb{Z}/2$-equivariant as well.
Proof. We start by denoting

\[ X(x, y) = X^1(x, y) \frac{\partial}{\partial x} + X^2(x, y) \frac{\partial}{\partial y} \]

Every smooth function \( h: T^2 \to \mathbb{R} \) has a Fourier series

\[ \mathcal{F}[h](x, y) = \sum_{k \in \mathbb{Z}^2} u_k \sin(k \cdot (x, y)) \]
\[ + \sum_{k \in \mathbb{Z}^2} v_k \cos(k \cdot (x, y)) \]

where the (real) coefficients \( u_k, v_k \) are given by

\[ u_k = \frac{1}{(2\pi)^2} \int_{T^2} h(x, y) \sin(k \cdot (x, y)) \, d(x, y) \]
\[ v_k = \frac{1}{(2\pi)^2} \int_{T^2} h(x, y) \cos(k \cdot (x, y)) \, d(x, y) \]

It is well known that if \( h \) is a smooth (\( C^\infty \)-) function, then the Fourier series \( \mathcal{F}[h] \) converges to \( h \) in any \( C^n \)-norm. (The corresponding statement is less neat if \( h \) is just a \( C^N \) function for some sufficiently big \( N \).)

If \( h \) is an odd function, meaning that we have \( h(-x, y) = -h(x, y) \) for all \((x, y) \in T^2\), then all the coefficients \( v_k \) are zero.

In what follows, we only suppose that \( X \) is a \( \mathbb{Z}/2 \)-equivariant vector field, since this is the case that is relevant to us here. The proof without this restriction is completely analogous.

So suppose the coefficients \( X^1(x, y) \) and \( X^2(x, y) \) of \( X \) have Fourier series

\[ X^i(x, y) = \sum_{k \in \mathbb{Z}^2} u_k^i \sin(k \cdot (x, y)) , \quad i = 1, 2. \]

We will write \( k = (k_1, k_2) \). By assumption, we have

\[ 0 \equiv \text{div}(X) = \sum_{k \in \mathbb{Z}^2} (k_1 u_k^1 + k_2 u_k^2) \cos(k \cdot (x, y)). \]

As \( u_{-k}^i = -u_k^i \), we must therefore have

\[ \begin{pmatrix} u_k^1 \\ u_k^2 \end{pmatrix} \cdot \begin{pmatrix} k_1 \\ k_2 \end{pmatrix} = 0 \]
for all \( k \in \mathbb{Z}^2 \). Every Fourier term \( W_k \) of the Fourier expansion \( X = \sum_k W_k \) of the divergence-free vector field \( X \) is therefore of the form

\[
W_k(x, y) = \sin(k \cdot (x, y)) \begin{pmatrix} u^1_k \\ u^2_k \end{pmatrix}.
\] (16)

In other words, \( W_k \) is a \( \mathbb{Z}/2 \) equivariant vector field of shearing type with integer shearing direction \( \mathbf{k} := (-k_2, k_1) \in \mathbb{Z}^2 \), as we may write the vector \( (u^1_k, u^2_k) \) as a scalar multiple of \( \mathbf{k} \). □

If \( X: T^2 \to \mathbb{R}^2 \) is a smooth vector field, it must be Lipschitz-continuous. In other words, there is a constant \( L > 0 \) such that

\[
\|X(x, y) - X(x', y')\| \leq L d((x, y), (x', y')),
\]

for all \( (x, y), (x', y') \in T^2 \), where \( d(-, -) \) is the Euclidean metric on \( T^2 \), and where \( \|\cdot\| \) denotes the Euclidean norm. We will find it convenient to denote by \( (x, y) \in \mathbb{R}^2 \) also some lift of the point on the torus. With this convention, the last inequality also implies

\[
\|X(x, y) - X(x', y')\| \leq L \|(x, y) - (x', y')\|.
\]

The proof of Theorem 3.3 above will essentially rely on the preceding Lemma, as well as on some classical results:

**Lemma 3.5** (Gronwall’s inequality). Let \( f, g: [a, b] \to \mathbb{R} \) be two continuous non-negative functions, and suppose that one has

\[
f(t) \leq A(t) + \int_a^t f(s) g(s) \, ds
\]

for all \( t \), for some non-negative increasing function \( A: [a, b] \to \mathbb{R} \). Then one has

\[
f(t) \leq A(t) \exp \left( \int_a^t g(s) \, ds \right)
\]

for all \( t \in [a, b] \).

This is Lemma 2 after Theorem 2.1.2 in [1], to which we refer for a proof.

**Lemma 3.6.** Let \( X: T^2 \to \mathbb{R}^2 \) be a vector field on \( T^2 \) with Lipschitz constant \( L \). We denote by \( \phi^t_X: T^2 \to T^2 \) its flow, satisfying

\[
\frac{d}{dt} \phi^t_X(p) = X(\phi^t_X(p)),
\]
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for all \( p \in T^2 \) and all \( t \in \mathbb{R} \). Then we have

\[
||\phi^t_X(p) - \phi^t_Y(q)|| \leq e^{Lt} ||p - q||
\]

(17)

for all \( p, q \in T^2 \), and for all \( t \geq 0 \).

For the proof we refer to [1, Lemma 3 after Theorem 2.1.2].

**Lemma 3.7.** If \( X \) and \( Y \) are two vector fields on \( T^2 \) with flows \( \phi^t_X \) and \( \phi^t_Y \). Then we have

\[
||\phi^t_X(p) - \phi^t_Y(p)|| \leq t ||X - Y||_{L^\infty(T^2)} e^{Lt},
\]

(18)

for all \( p \in T^2 \), and all \( t \geq 0 \). Here \( L \) is a Lipschitz constant for the vector field \( X \) or for \( Y \), and \( ||X||_{\infty} = \max_{p \in T^2} ||X(p)|| \) denotes the maximum (or supremum) norm.

**Proof.** Integrating the flow \( \phi^t_X \) of the vector field \( X \), we obtain the integral equation

\[
\phi^t_X(p) = p + \int_0^t X(\phi^s_X(p)) \, ds.
\]

This yields

\[
||\phi^t_X(p) - \phi^t_Y(p)|| \leq ||\int_0^t (X(\phi^s_X(p)) - Y(\phi^s_X(p)) + Y(\phi^s_X(p)) - Y(\phi^s_Y(p))) \, ds||
\]

\[
\leq t ||X - Y||_{L^\infty(T^2)} + L \int_0^t ||\phi^s_X(p) - \phi^s_Y(p)|| \, ds,
\]

where in the second inequality we have used the Lipschitz property of the vector field \( Y \), with Lipschitz constant \( L \). Hence the function \( f(t) := ||\phi^t_X(p) - \phi^t_Y(p)|| \) satisfies the integral inequality

\[
f(t) \leq t ||X - Y||_{\infty} + L \int_0^t f(s) \, ds.
\]

(19)

This implies \( f(t) \leq t ||X - Y||_{\infty} \exp(Lt) \) for all \( t \geq 0 \) by Gronwall’s inequality. The role of \( X \) and \( Y \) can be interchanged in the proof, so the Lipschitz constant can be either a suitable one for \( X \) or for \( Y \). \( \square \)

We also state an immediate combination of the preceding two Lemmas, using the triangle inequality.
Lemma 3.8. Let $X, Y$ be two vector fields on the torus, and let $p, q \in T^2$. Let $L$ be a Lipschitz constant either for $X$ or for $Y$, as above. Then we have
\[ \| \phi^t_X(p) - \phi^t_Y(q) \| \leq \| p - q \| e^{L t} + t \| X - Y \|_{\infty} e^{L t} \]
for all $p, q \in T^2$, and for all $t \geq 0$.

Notice that an isotopy $(\psi_t)_{t\in[0,1]}$ as in the statement of Theorem 3.3 gives rise to a time-dependent vector field $X_t$ which satisfies the differential equation
\[ \frac{d\psi^t(p)}{dt} = X_t(\psi^t(p)) \]
for any $p \in T^2$, and any $t \in [0,1]$. The end point we reach from $p \in T^2$ when integrating the time-dependent vector field from time $t_0$ to time $t$ is given by $\psi^t((\psi^{t_0})^{-1}(p))$.

Proof of Theorem 3.3. Let $\varepsilon > 0$ be given. The proof of Theorem 3.3 is an approximation in three steps.

(1) (First step) We approximate the ‘flow’ $\psi_t$ of the time-dependent vector field $X_t$ by a composition of flows of the finitely many time-independent vector fields given by $X_t$ taken at times $\frac{i}{n}$:
\[ X_i := X_{i/n}, \quad \text{for } i = 0, \ldots, n - 1. \]

More precisely, we define a continuous family of isotopies $(\Theta^t_{(X_j)})$
\[ \Theta^t_{(X_j)}: T^2 \to T^2 \]
in the following way: For $\frac{i}{n} \leq t \leq \frac{i+1}{n}$, we denote
\[ \Theta^t_{(X_j)} := \phi^t_{X_i} \circ \phi^1_{X_{i-1}} \circ \cdots \circ \phi^1_{X_0}. \]
(21)

In other words, we first flow along the vector field $X_0$ over a period of time $\frac{1}{n}$, then we flow along $X_1$ over a period of time $\frac{1}{n}$, and finally we flow along $X_i$ over a period of time $t - \frac{i}{n}$. By keeping track of the accumulated error in each time interval of length $\frac{1}{n}$, we will show that
\[ \| \Theta^t_{(X_j)} - \psi^t \|_{\infty} < \frac{\varepsilon}{3} \]
for all $t \in [0,1]$ if we choose $n$ large enough.
(2) (Second step) Each of the divergence free vector fields $X_i$ can be $C^0$-approximated by a vector field $Z_i$ which is a finite Fourier series, by Lemma 3.4 above. We define $\Theta^t(Z_i)$ to be the map defined completely analogously to the one defined in (21) above, but with the vector fields $Z_i$ instead of the $X_i$.

Again, by keeping track of the accumulated error, we will find that

$$\|\Theta^t(X_j) - \Theta^t(Z_j)\|_\infty < \frac{\varepsilon}{3}$$

for all $t \in [0, 1]$, if we choose each $Z_j$ sufficiently close to $X_j$.

(3) (Third step)

Each of the vector fields $Z_j$ is a finite Fourier series

$$Z_j = \sum_{r=0}^{m_j-1} W_r^{(j)}$$

where each $W_r^{(j)}$ is a Fourier term, and hence a shearing vector field in a direction in $\mathbb{Z}^2$ by Lemma 3.4 above. The map $\Theta^t(Z_j)$ of the second step is made out of a composition of flows $\phi_{Z_j}^t$ for some short period of time $t$ smaller than $\frac{1}{n}$. We finally will approximate the flow along $Z_j$ by successive flows along the summands $W_r^{(j)}$ in the following way.

For simplicity we omit the index $j$ from the notation in the vector field $Z_j$ and $W_r^{(j)}$ for a moment. We fix some $k \in \mathbb{N}$ (that also depends on $j = 0, \ldots, n - 1$). Suppose we have

$$\frac{1}{n} \left( \frac{i}{k} + \frac{r}{km} \right) \leq t \leq \frac{1}{n} \left( \frac{i + 1}{k} + \frac{r + 1}{km} \right),$$

for some $i = 0, \ldots, k - 1$, and some $r = 0, \ldots, m - 1$, and in particular $0 \leq t \leq \frac{1}{n}$, where $n$ was fixed in the first step. Then we define

$$\Xi^t(W_r) := \phi_{W_r}^{m(t - \frac{im + r}{kmm})} \circ \phi_{W_r^{m-1}}^{\frac{1}{km}} \circ \cdots \circ \phi_{W_r}^{\frac{1}{m}} \circ \phi_{W_0}^{\frac{1}{m}},$$

(22)

In other words, we first flow along $W_0$ during a period of time $\frac{1}{kmm}$, but with speed $m$, then along $W_1$ during a period of time $\frac{1}{kmm}$ with speed $m$, and so on. After we have flown along $W_{m-1}$, we repeat this process, flowing again along $W_0$ during a period of time $\frac{1}{kmm}$ with
speed \( m \), then along \( W_1 \) etc. In other words, \( \Xi(W_r) \) is a repeated iteration of flows along the vector fields \( W_0, \ldots, W_{m-1} \).

We will finally compare \( \Theta_t(Z_j) \) to the map \( \Omega_t(Z_j) \) which we define as follows. For \( \frac{i}{n} \leq t \leq \frac{i+1}{n} \), we denote with slight abuse of notation by

\[
\Omega_t(Z_j) := \Xi_{t-i/n}(W_i^{(i)}) \circ \Xi_{t-i/n-1}^{1/n}(W_i^{(i-1)}) \circ \cdots \circ \Xi_{t-i/n-1}^{1/n}(W_i^{(0)}) ,
\]

for \( i = 0, \ldots, n-1 \). We will show that if for each sum \( Z_i = \sum_{r=0}^{m-1} W_r^{(i)} \) we choose the ‘fineness’ \( k_i \) in the definition of \( \Xi(W_r^{(i)}) \) large enough, we will have

\[
\| \Theta_t(Z_j) - \Omega_t(Z_j) \|_\infty < \frac{\varepsilon}{3}
\]

for all \( t \in [0, 1] \).

3.1. **First step.** Let \( L > 0 \) be a Lipschitz constant for the time dependent vector fields \( X_t, t \in [0, 1] \). In other words, we assume that

\[
\| X_t(p) - X_t(q) \| \leq L \| p - q \|
\]

for all \( p, q \in T \), and all \( t \in [0, 1] \). The existence of such an \( L \) follows from a standard compactness argument when we consider the time-dependent vector field \( X_t \) as a map \( T^2 \times [0, 1] \to \mathbb{R}^2 \). It fact, it turns out useful to consider the time-independent vector fields \( X_i \) as maps \( T^2 \times [0, 1] \to \mathbb{R}^2 \) which are constant in the first factor.

As in the proof of Lemma 3.7, we integrate the vector fields \( X_0 \) as well as the time-dependent vector field \( X_t \) to obtain

\[
\| \psi^t(p) - \phi^t_{X_0}(p) \| \leq \int_0^t \| (X_s(\psi^s(p)) - X_0(\phi^s_{X_0}(p))) \| \, ds
\]

\[
\leq \int_0^t (L \| \psi^s(p) - \phi^s_{X_0}(p) \| + \| X_s - X_0 \|_{L^\infty(T^2 \times [0,t])}) \, ds .
\]

Gronwall’s inequality now implies

\[
\| \psi^t(p) - \phi^t_{X_0}(p) \| \leq t \| X_s - X_0 \|_{L^\infty([0,t] \times T^2)} e^{Lt}
\]

for all \( t \geq 0 \), and all \( p \in T^2 \).

Similarly, for some integer \( i \geq 0 \), we can integrate both the time-dependent vector field \( X_t \) and \( X_i = X_{i/n} \) starting at time \( \frac{i}{n} \), and the same argument gives

\[
\| \psi^t((\psi_{\frac{i}{n}})^{-1}(p)) - \phi^t_{X_i}(p) \| \leq (t - \frac{i}{n}) \| X_s - X_i \|_{L^\infty(T \times [\frac{i}{n}, t])} e^{L(t-\frac{i}{n})}
\]
for all \( t \geq \frac{i}{n} \), and all \( p \in T^2 \).

However, at time \( \frac{i}{n} \) we will in general already have taken up an error between flowing along the time-dependent vector field, and between flowing first along \( X_0 \) during time \( \frac{i}{n} \), then along \( X_1 \) during time \( \frac{i}{n} \), and so on. What we therefore need is the following inequality that we simply obtain from (25) and Lemma 3.8 using the triangle inequality:

\[
\| \psi^t((\psi^{\frac{i}{n}})^{-1}(p)) - \phi^t_{X_i}(q) \| \leq (t - \frac{i}{n}) \| X_s - X_i \|_{L^\infty(T \times [\frac{i}{n}, \frac{i}{n}+1])} e^{L(t - \frac{i}{n})} + \| p - q \| e^{L(t - \frac{i}{n})} ,
\]

(26)

for all \( t \geq \frac{i}{n} \), and all \( p, q \in T^2 \). We are now able to estimate the accumulated error.

**Lemma 3.9.** For \( \Theta^t(X_j) \) as above, and for any \( k = 0, \ldots, n-1 \), we have

\[
\| \psi^t(p) - \Theta^t(X_j)(p) \| \leq \frac{1}{n} \sum_{i=0}^{k-1} \| X_s - X_i \|_{L^\infty([\frac{i}{n}, \frac{i}{n}+1] \times T^2)} e^{L(t - \frac{i}{n})} + (t - \frac{k}{n}) \| X_s - X_k \|_{L^\infty([\frac{k}{n}, \frac{k}{n}+1] \times T^2)} e^{L(t - \frac{k}{n})} ,
\]

(27)

for all \( t \geq \frac{k}{n} \) and all \( p \in T^2 \). In particular, we have

\[
\| \psi^t(p) - \Theta^t(X_j)(p) \| \leq \frac{1}{n} \sum_{i=0}^{n-1} \| X_s - X_i \|_{L^\infty([\frac{i}{n}, \frac{i}{n}+1] \times T^2)} e^L
\]

for all \( t \in [0, 1] \) and all \( p \in T^2 \).

**Proof.** We proof the claim by induction on \( k \). For \( k = 0 \) the claim follows from what we have established in (24) above.

Suppose it is true for some \( k \leq n - 1 \), and assume we have \( t \geq \frac{k+1}{n} \). By inequality (26) we have

\[
\| \psi^t(p) - \Theta^t(X_j)(p) \| \leq \| \psi^t((\psi^{\frac{k+1}{n}})^{-1}(\psi^{\frac{k+1}{n}}(p))) - \phi^t_{X_{k+1}}(\Theta^{\frac{k+1}{n}}(p)) \|
\]

\[
\leq (t - \frac{k+1}{n}) \| X_s - X_0 \|_{L^\infty(T \times [\frac{k+1}{n}, \frac{k+1}{n}+1])} e^{L(t - \frac{k+1}{n})} + \| \psi^{\frac{k+1}{n}}(p) - \Theta^{\frac{k+1}{n}}(p) \| e^{L(t - \frac{k+1}{n})} .
\]
We can now use our induction hypothesis for the last term, which yields
\[
\left\| \psi^t(p) - \Theta^t_{(X_j)}(p) \right\| \leq (t - \frac{k + 1}{n}) \left\| X_s - X_{k+1} \right\|_{L^\infty(T \times [\frac{k+1}{n}, t])} e^{L(t - \frac{k+1}{n})} \\
+ \left( \frac{1}{n} \sum_{i=0}^{k} \left\| X_s - X_i \right\|_{L^\infty([\frac{i+1}{n}, \frac{i+1}{n}] \times T^2)} e^{L(t - \frac{k+1}{n} - \frac{1}{n})} \right) e^{L(t - \frac{k+1}{n})}.
\]
This simplifies to the statement we wished to show for \( k + 1 \).

Finally, notice that the time-dependent vector field \( X_t \), when seen as a map \( T^2 \times [0, 1] \to \mathbb{R}^2 \), is equicontinuous. In particular, there exists some \( \delta > 0 \) such that whenever \( |s - t| < \delta \), we have
\[
\left\| X_s(p) - X_t(p) \right\| \leq \frac{\varepsilon}{3} e^{L}
\] (28)
for all \( p \in T^2 \).

The desired first estimate now follows quickly. We choose the integer \( n \) so large that \( \frac{1}{n} < \delta \). Then we will have
\[
\left\| X_s - X_i \right\|_{L^\infty([\frac{i+1}{n}, \frac{i+1}{n}] \times T^2)} \leq \frac{\varepsilon}{3} e^{L}
\]
for \( i = 0, \ldots, n - 1 \), and hence by the Lemma 3.9 we see that
\[
\left\| \psi^t - \Theta^t_{(X_j)} \right\|_{\infty} < \frac{\varepsilon}{3}
\]
for all \( t \in [0, 1] \).

3.2. Second step.

**Lemma 3.10.** For \( \Theta^t_{(X_j)} \) and \( \Theta^t_{(Z_j)} \) as above, and for any \( k = 0, \ldots, n - 1 \), we have
\[
\left\| \Theta^t_{(X_j)}(p) - \Theta^t_{(Z_j)}(p) \right\| \leq \frac{1}{n} \sum_{i=0}^{k-1} \left\| X_i - Z_i \right\|_{\infty} e^{L(t - \frac{k}{n})} \\
+ \left( t - \frac{k}{n} \right) \left\| X_k - Z_k \right\|_{\infty} e^{L(t - \frac{k}{n})}
\]
for all \( t \geq \frac{k}{n} \) and all \( p \in T^2 \). In particular, we have
\[
\left\| \Theta^t_{(X_j)}(p) - \Theta^t_{(Z_j)}(p) \right\| \leq \frac{1}{n} \sum_{i=0}^{n-1} \left\| X_i - Z_i \right\|_{\infty} e^{L}.
\]
for all $t \in [0, 1]$ and all $p \in T^2$. Here $L$ is the same Lipschitz constant that worked for the time-dependent vector field $X_t$, as above.

**Proof.** This is a consequence of Lemma 3.7 above, together with an induction argument analogous to the one of the Proof of Lemma 3.9. □

In order to establish the second step approximation, we choose $Z_i$ so that

$$\|X_i - Z_i\|_\infty \leq \frac{\varepsilon}{3 e^L}$$

for each $i = 0, \ldots, n - 1$. Lemma 3.10 now implies that we have

$$\|\Theta^t_{(X_j)} - \Theta^t_{(Z_j)}\|_\infty < \frac{\varepsilon}{3}$$

for all $t \in [0, 1]$.

### 3.3. Third step.

We will need the following Lemma, which at least in weaker formulations seems to be a classical result.

**Lemma 3.11.** Suppose we have a finite sum of vector fields $Z = W_1 + \cdots + W_m$ on $T^2$. Then there is a constant $C$, depending only on the vector fields $W_1, \ldots, W_m$, such that we have

$$\|\phi^t_Z - \phi^t_{W_m} \circ \cdots \circ \phi^t_{W_1}\|_\infty \leq \frac{t^2}{2} C e^{Lt},$$

for all $0 \leq t \leq 1$, and where $L$ is a Lipschitz constant for the vector field $Z$.

**Proof.** For simplicity, we will give a proof of this result in the case $m = 2$. It will be clear that the more general result is completely analogous.

We recall the definition (or one of various equivalent ones) of the Lie-bracket of two vector fields $X$ and $Y$,

$$[X, Y](p) = \lim_{h \to 0} \frac{\phi^h_Y)_*X(\phi^{-h}_Y(p)) - X(p)}{h} = \frac{d}{dt}[(\phi^t_Y)_*(X)](p)|_{t=0},$$

where $(\phi^t_Y)_*$ denotes the derivative of the map $\phi^t_Y$. Stated slightly differently, we can say that

\begin{equation}
(\phi^t_Y)_*X(p) = X(\phi^t_Y(p)) + t[X, Y](\phi^t_Y(p)) + R_{X,Y}(t, \phi^t_Y(p)),
\end{equation}

where the ‘error’ $R_{X,Y}$ satisfies

$$\lim_{t \to 0} \frac{R_{X,Y}(p, t)}{t} = 0$$
and this convergence is uniform in $p$, as we are on a compact manifold.

Using this, we start by differentiating the composite of the flows along $W_1$ and $W_2$ and obtain

$$
\frac{d}{dt} \phi_{W_2}^t(\phi_{W_1}^t(p)) = \frac{d\phi_{W_2}^t(\phi_{W_1}^t(p))}{dt} + (\phi_{W_2}^t)_*(\frac{d\phi_{W_1}^t(p)}{dt})
= W_2(\phi_{W_2}^t(\phi_{W_1}^t(p))) + (\phi_{W_2}^t)_*W_1(\phi_{W_1}^t(p))
= W_2(\phi_{W_2}^t(\phi_{W_1}^t(p))) + W_1(\phi_{W_2}^t(\phi_{W_1}^t(p)))
+ t[W_1, W_2](\phi_{W_2}^t(\phi_{W_1}^t(p))) + R_{W_1, W_2}(p, t),
$$

where we have made use of (29). We can integrate this equation and compare this to the flow $\phi_{W_1 + W_2}^t$ of the vector field $W_1 + W_2$, obtaining

$$
\|\phi_{W_1 + W_2}^t(p) - \phi_{W_2}^t(\phi_{W_1}^t(p))\|
\leq \int_0^t \|(W_1 + W_2)(\phi_{W_1 + W_2}^s(p)) - (W_1 + W_2)\phi_{W_2}^s(\phi_{W_1}^s(p))\| \, ds
+ \int_0^t s([W_1, W_2](\phi_{W_2}^s(\phi_{W_1}^s(p)))) + \frac{1}{s} R_{W_1, W_2}(\phi_{W_2}^s(\phi_{W_1}^s(p)), s)\| \, ds
\leq L \int_0^t \|\phi_{W_1 + W_2}^s(p) - \phi_{W_2}^s(\phi_{W_1}^s(p))\| \, ds + \frac{t^2}{2} C_{W_1, W_2},
$$

where $L$ is a Lipschitz constant for the vector field $W_1 + W_2$, and where $C_{W_1, W_2}$ is a constant that can be taken to be

$$
C_{W_1, W_2} = \|[W_1, W_2]\|_{L^\infty(T^2)} + \|R_{W_1, W_2}/s\|_{L^\infty(T^2 \times [0,1])}.
$$

Using Gronwall’s inequality, this yields

$$
\|\phi_{W_1 + W_2}^t - \phi_{W_2}^t \circ \phi_{W_1}^t\|_{L^\infty} \leq \frac{t^2}{2} C_{W_1, W_2} e^{Lt}
$$

for all $0 \leq t \leq 1$. The general case follows similarly, only the constant $C_{W_1, \ldots, W_m}$ will contain a more complicated expression in terms of commutators, and commutators of commutators etc. $\Box$

Using the triangle inequality and Lemma 3.11, we obtain

$$
\|\phi_Z^t(p) - \phi_{W_{m-1}}^t \circ \cdots \circ \phi_{W_0}^t(q)\| \leq \|p - q\| e^{Lt} + \frac{t^2}{2} C e^{Lt},
$$

for all $p, q \in T^2$, where $L$ and $Z$ are as in the preceding Lemma.
We now have to compare the flows of the vector fields $Z_j$ to the repeated iteration of flows along the vector fields $W_0^{(j)}, \ldots, W_{m_j-1}^{(j)}$ occurring in the definition of $\Xi^t$ defined in (22) above. This is easier at the time intervals which are multiples of $\frac{1}{nk_j}$, and we need a more coarse estimate at the times in between.

**Lemma 3.12.** Suppose we have a finite sum of vector fields $Z = W_1 + \cdots + W_m$ on $T^2$, and we assume $L$ is a Lipschitz constant for the vector field $Z$. Let $C$ be a constant assureing the conclusion of Lemma 3.11 above.

(i) For $i = 0, \ldots, k - 1$, we have the estimate

$$\|\phi_{\frac{t}{kn}}(p) - \Xi_{(W_r)}^i(q)\| \leq \|p - q\| e^{L \frac{t}{kn}} + i \frac{(1/kn)^2}{2} Ce^{L \frac{t}{kn}},$$

for any $p, q \in T^2$.

(ii) For $i \leq t \leq i + \frac{1}{kn}$ we have

$$\|\phi_{\frac{t}{kn}}(p) - \Xi_{(W_r)}^i(q)\| \leq \|p - q\| e^{L \frac{t}{kn}} + i \frac{(1/kn)^2}{2} Ce^{L \frac{t}{kn}} + (t - \frac{i}{kn})(\|Z\|_\infty + \|W_{m-1}\|_\infty + \cdots + \|W_{0}\|_\infty).$$

*(33)*

**Proof.** The first part is easily shown by induction on $i$, using the estimate (32) after the preceding Lemma.

To obtain the second part, for $i \leq t \leq i + \frac{1}{kn}$, we compare this to the nearby situation at time $\frac{i}{kn}$, and we obtain

$$\|\phi_{\frac{t}{kn}}(p) - \Xi_{(W_r)}^i(q)\| \leq \|\phi_{\frac{t}{kn}}(p) - \phi_{\frac{i}{kn}}^t(p)\|$$

$$+ \|\phi_{\frac{i}{kn}}^t(p) - \Xi_{(W_r)}^i(q)\| + \|\Xi_{(W_r)}^i(q) - \Xi_{(W_r)}^t(q)\|$$

with the triangle inequality. The first and the third term of the right hand side are easily estimated by the inequality

$$\|\phi_{t}^t(p) - p\| \leq t\|X\|_\infty$$

that holds for the flow of any vector field $X$, for $t \geq 0$. The second term is estimated using the first part of the Lemma. All three estimates together yield the desired result. \(\square\)
We can now establish our final estimate. We recall that we have a decomposition $Z_j = W_{0}^{(j)} + \cdots + W_{m_j-1}^{(j)}$ for any $j = 0, \ldots, n-1$. The constant $k$ chosen in the decomposition of the time interval $[0, 1/n]$ into intervals of length $1/kn$ also depends on $j$, hence we denote these numbers by $k_j$ and so we do with the constants $C_j$ of Lemma 3.11 above. Notice that since $Z_j$ is taken to be a finite sum in the Fourier series of Lemma 3.4 above, and this series converges in any $C^m$-topology, we may assume that the Lipschitz constant $L$ for $X_j$ also works for the $Z_j$.

**Lemma 3.13.** For $\Theta_t^{(Z_j)}$ and $\Omega_t^{(Z_j)}$ as above, and for any $l = 0, \ldots, n-1$, we have

$$\|\Theta_t^{(Z_j)}(p) - \Omega_t^{(Z_j)}(p)\| \leq \sum_{j=0}^{l-1} \left( \frac{1}{k_j n} (\|Z_j\|_{\infty} + \|W_{m_j-1}^{(j)}\|_{\infty} + \cdots + \|W_0^{(j)}\|) + \frac{1}{k_j n^2} C_j e^{L \frac{1}{k_j n}} \right) e^{L(t - \frac{j}{n})}$$

for all $\frac{j}{n} \leq t \leq \frac{j+1}{n}$, and for all $p \in T^2$. In particular, we have

$$\|\Theta_t^{(X_j)}(p) - \Omega_t^{(Z_j)}(p)\| \leq \sum_{j=0}^{n-1} \left( \frac{1}{k_j n} (\|Z_j\|_{\infty} + \|W_{m_j-1}^{(j)}\|_{\infty} + \cdots + \|W_0^{(j)}\|) + \frac{1}{k_j n^2} C_j e^{L \frac{1}{k_j n}} \right) e^{L}$$

(34)

for all $t \in [0, 1]$ and all $p \in T^2$. Here $L$ may be assumed to be the same Lipschitz constant that worked for the time-dependent vector field $X_t$, as above.

**Proof.** This is a consequence of Lemma 3.13 above, together with an induction argument analogous to the one of the Proof of Lemma 3.9 $\Box$

We can now finish the proof of the third step. We simply choose the $k_j$ which give the number of equidistant subdivisions of $[0, \frac{1}{n}]$ to be so large that

$$\frac{1}{k_j} (\|Z_j\|_{\infty} + \|W_{m_j-1}^{(j)}\|_{\infty} + \cdots + \|W_0^{(j)}\| + \frac{1}{n} C_j e^{L}) e^{L} \leq \frac{\varepsilon}{3}$$
for each $j = 0, \ldots, n - 1$. Lemma 3.13 now implies that we have

$$\|\Theta^t_{(Z_j)} - \Omega^t_{(Z_j)}\|_\infty < \frac{\varepsilon}{3}$$

for all $t \in [0, 1]$.

We finally notice that $(\Omega^t)$ is a finite composition of shearing maps for any $t \in [0, 1]$. In fact, a vector field which is a Fourier term in Lemma 3.4 above has the form

$$W_k(x, y) = \sin(k \cdot (x, y)) \begin{pmatrix} u^1_k \\ u^2_k \end{pmatrix},$$

where $k = (k_1, k_2) \in \mathbb{Z}^2$, and $(u^1_k, u^2_k) \in \mathbb{R}^2$ is a multiple of $\overline{k} = (-k_2, k_1) \in \mathbb{Z}^2$, see [16] above. The flow $\phi^t_W$ takes the form

$$\phi^t_{W_k}(x, y) = (x, y) + t W_k(x, y).$$

This is a shearing map in a direction of $\mathbb{Z}^2$.

Therefore, for any $t \in [0, 1]$, the maps $\Xi^t_{(Z_j)}$, and hence also the maps $\Omega^t_{(Z_j)}$, are compositions of finitely many shearing maps by construction.

As for the $\mathbb{Z}/2$-equivariance, we simply notice that all maps and vector fields in the construction are $\mathbb{Z}/2$-equivariant, and hence so is $\Omega^t_{(Z_j)}$. This terminates the proof of Theorem 3.3 if we define the desired map $\phi^t$ by $\phi^t := \Omega^t_{(Z_j)}$.

**Remark 3.14.** There is a straightforward generalisation of the preceding theorem and its proof to isotopies of the $n$-dimensional torus $T^n$ which preserve the $n$-dimensional Euclidean volume form, and in particular, to isotopies through symplectomorphisms in even dimensions. However, as we do not need this more general statement, we have contented ourselves to the 2-dimensional case.

4. **Isotopies of the pillowcase approximated by holonomy perturbations, main technical result**

The results of Proposition 2.1 and Theorem 3.3 will now yield our main technical result.

Recall that the representation variety of the 2-dimensional torus $R(T^2)$ is a 2-dimensional sphere. Let us denote by $\widehat{R}(T^2)$ the double cover branched
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over the four singular points of \( R(T^2) \). Working with diffeomorphisms, isotopies etc. of \( R(T^2) \), which fix the four singular points, is equivalent to working with diffeomorphisms, isotopies etc. of \( \hat{R}(T^2) \), which are \( \mathbb{Z}/2 \)-equivariant and fix the four fixed points there. There is an identification

\[
\mathbb{R}^2/2\pi\mathbb{Z}^2 \to \hat{R}(T^2) \subseteq \text{Hom}(\mathbb{Z}^2, SU(2))
\]

\[(\alpha, \beta) \mapsto \rho \text{ with } \rho(m) = \begin{bmatrix} e^{i\alpha} & 0 \\ 0 & e^{-i\alpha} \end{bmatrix} \text{ and } \rho(l) = \begin{bmatrix} e^{i\beta} & 0 \\ 0 & e^{-i\beta} \end{bmatrix}.
\]

We will think of a geometry on \( \hat{R}(T^2) \) which is the one induced from the canonical Euclidean structure on \( \mathbb{R}^2/2\pi\mathbb{Z}^2 \) via this identification.

In the following theorem, we will consider the 3-manifold \( M = [0, 1] \times S^1 \times S^1 \), a thickened torus. Let \( \mu \) denote a 2-form on \([0, 1] \times S^1 \times S^1\) with integral 1 and with compact support in the interior. If \( \iota_k : [0, 1] \times S^1 \times S^1 \to M \) is an embedding, then we denote with \( \mu_k \) the 2-form on \( M \) which when pulled back by \( \iota_k \) is the 2-form on \([0, 1] \times S^1 \times S^1\) which is obtained by pulling back \( \mu \) from \([0, 1] \times S^1 \times S^1\) to \([0, 1] \times S^1 \times S^1\) via the projection.

Let \( z_0 = (x_0, y_0) \in S^1 \times S^1 \) be some chosen base point. For \( i = 0, 1 \) we will consider the closed curves

\[
m_i = \{r\} \times S^1 \times \{y_0\} \quad \text{and} \quad l_i = \{r\} \times \{x_0\} \times S^1.
\]

(35) on the boundary of \( M \).

**Definition 4.1.** Let the function \( \theta^t : \mathcal{A} \to \Omega^2(Y; \mathfrak{su}(E)) \) be defined in the following way. For \( \frac{k}{n} \leq t \leq \frac{k+1}{n} \) and some \( k = 0, \ldots, n-1 \),

\[
\theta^t_{\{t_k, \chi_k\}}(A) := n(t - \frac{k}{n}) \cdot \chi_k(\text{Hol}_{t_k}(A)) \mu_k + \sum_{l=0}^{k-1} \chi_l(\text{Hol}_{t_l}(A)) \mu_l.
\]

For any \( t \), this corresponds to some holonomy perturbation of the flatness equation by some holonomy perturbation data. The function interpolates between \( \theta^0 = 0 \) and the full holonomy perturbation determined by \( \{t_k, \chi_k\}_{k=0}^{n-1} \) for \( t = 1 \).

For any \( t \in [0, 1] \), we denote by \( R(M; \theta^t_{\{t_k, \chi_k\}}) \) the equivalence classes of connections \( A \in \mathcal{A} \) which solve the equation

\[
F_A = \theta^t_{\{t_k, \chi_k\}}(A),
\]

(36) and we call this the one-parameter family of perturbed flat connections determined by the holonomy perturbation data \( \{t_k, \chi_k\} \).
The following Theorem summarises the results of Sections 2 and 3.

**Theorem 4.2.** Let $M = [0, 1] \times S^1 \times S^1$ be a thickened torus, and let $E \to M$ be the trivial $SU(2)$-bundle over $M$. Let

$$
\psi: [0, 1] \times \hat{R}(T^2) \to \hat{R}(T^2)
$$

$$(t, (\alpha, \beta)) \mapsto \psi_t(\alpha, \beta)
$$

be a smooth map which is a $\mathbb{Z}/2$-equivariant isotopy through area-preserving maps with $\psi_0 = id$.

Then for any $\varepsilon > 0$, there is a $\mathbb{Z}/2$-equivariant isotopy through area-preserving maps

$$
\phi: [0, 1] \times \hat{R}(T^2) \to \hat{R}(T^2)
$$

$$(t, (\alpha, \beta)) \mapsto \phi_t(\alpha, \beta)
$$

which is continuous in $t$, smooth in $(\alpha, \beta)$, such that $\phi_0 = id$, and such that the following holds:

1. $\phi_t$ is $\varepsilon$-close to $\psi_t$ for all $t \in [0, 1]$,

$$
d(\psi_t(\alpha, \beta), \phi_t(\alpha, \beta)) < \varepsilon
$$

for all $t \in [0, 1]$ and all $(\alpha, \beta) \in \hat{R}(T^2)$, where $d$ denotes the metric induced by the Euclidean structure.

2. There is some $n \in \mathbb{N}$, and for $k = 0, \ldots, n - 1$ there are class functions $\chi_k: SU(2) \to \mathbb{R}$ corresponding to smooth even, $2\pi$-periodic functions $g_k: \mathbb{R} \to \mathbb{R}$ determined by

$$
\chi_k \left( \begin{bmatrix} e^{it} & 0 \\ 0 & e^{-it} \end{bmatrix} \right) = g_k(t),
$$

there are matrices

$$
A_k = \begin{pmatrix} a_k & c_k \\ b_k & d_k \end{pmatrix} \in SL(2, \mathbb{Z}),
$$

such that for $\frac{k}{n} \leq t \leq \frac{k+1}{n}$ we have

$$
\phi_t = \zeta_k^{n(t-\frac{k}{n})} \circ \zeta_{k-1}^1 \circ \cdots \circ \zeta_0^1
$$

for any $k = 0, \ldots, n - 1$.

Here the maps $\zeta_k^s$ are defined for $s \in [0, 1]$, and for $k = 0, \ldots, n - 1$ by the equation

$$
\zeta_k^s = A_k \circ \chi_k^s \circ A_k^{-1},
$$
where
\[
\chi^s_{f_k}: \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \mapsto \begin{pmatrix} \alpha + s \cdot f_k(\beta) \\ \beta \end{pmatrix},
\]
and where \( f_k \) is the derivative of \( g_k \).

(3) The nested sequence of embeddings
\[
\iota_k: [0, 1] \times S^1 \times S^1 \to M
\]
\[
\begin{pmatrix} s, \begin{pmatrix} z \\ w \end{pmatrix} \end{pmatrix} \mapsto \begin{pmatrix} k + s/n, \begin{pmatrix} a_k & c_k \\ b_k & d_k \end{pmatrix} \begin{pmatrix} z \\ w \end{pmatrix} \end{pmatrix},
\]
for \( k = 0, \ldots, n - 1 \) define holonomy perturbation data \( \{\iota_k, \chi_k\} \) with the following significance.

Suppose \( A \) solves the perturbed flatness equation (36) for some \( t \in [0, 1] \). Then \( A \) is reducible, and up to a gauge-transformation we may suppose that \( A \) respects the splitting \( E = M \times (\mathbb{C} \oplus \mathbb{C}) \). If we write the holonomies along the curves \( m_r, l_r \) for \( r = 0, 1 \) at the boundary as
\[
\text{Hol}_{m_r}(A) = \begin{bmatrix} e^{i\alpha_r} & 0 \\ 0 & e^{-i\alpha_r} \end{bmatrix}, \quad \text{and} \quad \text{Hol}_{l_r}(A) = \begin{bmatrix} e^{i\beta_r} & 0 \\ 0 & e^{-i\beta_r} \end{bmatrix},
\]
then we have
\[
\begin{pmatrix} \alpha_1 \\ \beta_1 \end{pmatrix} = \phi_t \begin{pmatrix} \alpha_0 \\ \beta_0 \end{pmatrix} \quad \text{or} \quad \begin{pmatrix} \alpha_1 \\ \beta_1 \end{pmatrix} = \phi_t \begin{pmatrix} -\alpha_0 \\ -\beta_0 \end{pmatrix},
\]
up to a composition in the \( t \)-variable with an increasing piecewise-linear homeomorphism \([0, 1] \to [0, 1]\).

A slightly more conceptual reformulation of the last statement is the following: For equivalence classes of connections in the one-parameter family of holonomy-perturbed representation varieties \( R(M; \theta^t_{\{\iota_k, \chi_k\}}) \) defined in 4.1 above we can take the holonomy at either boundary side of \( M \) to obtain restriction maps
\[
r_{\pm}: R(M; \theta^t_{\{\iota_k, \chi_k\}}) \to R(T^2),
\]
where \( r_- \) is the restriction to \( \{0\} \times S^1 \times S^1 \), and \( r_+ \) is the restriction to \( \{1\} \times S^1 \times S^1 \). Then the following diagram commutes:

\[
\begin{array}{ccc}
R(M; \theta^t_{\{\iota_k, \chi_k\}}) & \\
\downarrow r_- & \nearrow \phi_t \downarrow r_+ & \\
R(T^2) & \rightarrow & R(T^2),
\end{array}
\]
where here $\phi_t$ denotes the map between the pillowcase induced by $\phi_t$.

Proof. This follows at once from Proposition 2.5 and Theorem 3.3. □

5. Holonomy Perturbations on the 0-surgery of a Knot

We will now turn to the more concrete situation we have in mind. The 3-manifold we will consider is $Y = Y_0(K)$, by which we denote the 0-framed surgery on a knot $K$ in $S^3$. To set up our notation, we denote by $N(K)$ a tubular neighbourhood of the knot $K$ diffeomorphic to a solid torus. Then $Y$ is obtained by glueing to $Y(K) = S^3 \setminus N(K)$ the solid torus $S^1 \times D^2$ in such a way that longitudes of $K$ are mapped to circles $\{pt\} \times \partial D^2$.

The bundle $E \to Y$ is chosen in such a way that its first Chern-class $c_1(E)$ has odd evaluation on a Seifert surface of $K$ in $Y(K)$ capped off to a closed surface with a disc of the solid torus glued to $Y(K)$. Equivalently, we require $c_1(E)$ to be the Poincaré dual to an odd multiple of a meridian of $K$ inside $Y(K)$. Such a meridian generates $H_1(Y; \mathbb{Z}) \approx \mathbb{Z}$.

The manifold $Y(K)$ has boundary a 2-torus, and it has the homology of a circle. Hence the restriction $E \to Y(K)$ is necessarily trivial, and we suppose we have a trivialization chosen under which $\theta$ corresponds to the trivial connection in the determinant line. In particular, the curvature $F_{\theta}$ of $\theta$ is trivial over $Y(K)$.

The unperturbed representation variety $R^w(Y)$ for the 0-framed surgery $Y$ on $K$ in our situation is closely related to the representation variety

$$R(K) = \text{Hom}(\pi_1(S^3 \setminus K), SU(2))/SU(2),$$

of homomorphisms of the knot group into $SU(2)$, up to conjugation. Equivalently, we can understand $R(K)$ as a space of flat $SU(2)$-connections in the trivial bundle over $S^3 \setminus K$, up to gauge transformations.

The following is a classical fact, relating $R^w(Y)$ to representations of $R(K)$ to which some boundary conditions are imposed.

Proposition 5.1. There is a homeomorphism induced by the holonomy

$$R^w(Y_0(K)) \cong \{ \rho \in R(K) \mid \rho(l) = (-1)^k \text{id} \},$$

where $l$ is a longitude of the knot $K$, and where $k$ is equal to the evaluation $\langle c_1(E), \hat{\Sigma}_{\text{Seif}} \rangle$ on a Seifert surface of the knot capped off with a disc from
the solid torus which is glued in to produce the 0-surgery. If $k$ is odd, then the representations in $R(K)$ extend as $SO(3)$-representations of $\pi_1(Y_0(K))$ which do not lift to $SU(2)$-representations, and if $k$ is even they do extend as $SU(2)$-representations.

Proof. This follows from the Proof of Proposition 5.3 below if we take only class functions which are constant to 0 in the choice of holonomy perturbations. □

There is a collar neighbourhood of $\partial Y(K)$ that we may identify with $M = [0, 1] \times S^1 \times S^1$, such that the boundary $\partial Y(K)$ corresponds to $\{0\} \times S^1 \times S^1$. Let $z_0 = (x_0, y_0) \in S^1 \times S^1$ be some chosen base point. We suppose also that the identification is chosen so that of the closed curves

$m_r = \{r\} \times S^1 \times \{y_0\}$ and

$l_r = \{r\} \times \{x_0\} \times S^1,$

on the boundary of $M$ the curves $m_r$ correspond to meridians of the knot $K$, and such that the curves $l_r$ correspond to longitudes of $K$, for $r = 0, 1$.

By assumption, our bundle $E$ is trivialised over $Y(K)$, and so in particular over the thickened torus $M$. Just as in Theorem 4.2 or equivalently in Section 2 we choose a family of nested holonomy perturbations $\{\iota_k, \chi_k\}$ in $M$.

There is a similar description for the perturbed representation variety in terms of the representation variety $R(K)$ with some given boundary condition, analogous to Proposition 5.1 as we shall describe next. Notice first that the 3-manifold $Y_0(K)$ is given as a union of three pieces,

$Y_0(K) = (S^3 \setminus (N(K) \cup M)) \cup M \cup S^1 \times D^2.$

The first piece is diffeomorphic to $S^3 \setminus N(K)$ and homotopy equivalent to the knot complement, the second is the thickened torus, and the last one is a solid torus. The holonomy perturbations only have support in $M$, and our bundle is set up in such a way that any solution $A \in \mathcal{A}$ of the perturbed flatness equations

$(F_A)_0 = \sum_{k=0}^{n-1} \chi'_k(\text{Hol}_{0_k}(A)) \mu_k,$
is flat over the piece $S^3 \setminus N(K) \setminus M^\circ$, and such that its curvature over $S^1 \times D^2$ is equal to $\frac{1}{2} F_\theta \cdot \text{id}$, where $\theta$ is the fixed connection in the determinant line bundle $w$ of the Hermitian bundle $E \to Y_0(K)$.

In particular, we obtain a restriction map

$$r: R^w_{\{\iota_k, \chi_k\}}(Y_0(K)) \to R(K)$$

$$[A] \mapsto [A|_{S^3 \setminus (N(K) \cup M^\circ)}] \quad (38)$$

We introduce some notation, following Kronheimer and Mrowka [32].

**Definition 5.2.** For a flat connection $[A] \in R(K)$, the holonomy around a meridian $m$ and a longitude $l$ of $K$ define elements $\pm (\alpha(A), \beta(A)) \in R(T^2)$ if we assume that $A$ is chosen in its conjugacy class such that

$$\text{Hol}_m(A) = \begin{bmatrix} e^{i\alpha(A)} & 0 \\ 0 & e^{-i\alpha(A)} \end{bmatrix}, \quad \text{and} \quad \text{Hol}_l(A) = \begin{bmatrix} e^{i\beta(A)} & 0 \\ 0 & e^{-i\beta(A)} \end{bmatrix}.$$ 

If $S \subseteq R(T^2)$ is some $\mathbb{Z}/2$-invariant subset, then we define

$$R(K|S) := \{ [A] \in R(K) \mid \pm (\alpha(A), \beta(A)) \in S \}.$$ 

In other words, elements in $R(K|S)$ are conjugacy classes of flat connections on $S^3 \setminus N(K) \setminus M^\circ$ such that the restriction to the boundary torus has holonomy in $R(T^2)$ required to lie in the set $S$.

**Proposition 5.3.** For the bundle $E \to Y_0(K)$ with determinant line bundle $w$ and fixed reference connection $\theta$ as chosen above, the restriction map $r: R^w_{\{\iota_k, \chi_k\}}(Y_0(K)) \to R(K)$ defined in (38) above yields a homeomorphism

$$R^w_{\{\iota_k, \chi_k\}}(Y) \cong R(K|\tilde{\phi}_1(C)),$$

where $C = \{ (\alpha, \beta) | \beta = \pi \} \subseteq R(T^2)$, and where $\tilde{\phi}_1: R(T^2) \to R(T^2)$ is the map of the pillowcase induced by the holonomy perturbation data $\{\iota_k, \chi_k\}$ according to Theorem [4.2] above.

**Proof.** Any connection $[A] \in R^w_{\{\iota_k, \chi_k\}}(Y)$ has curvature $F_A = \frac{1}{2} F_\theta \cdot \text{id}$ over $S^1 \times D^2$. By Chern-Weil theory we have

$$\int_{\hat{\Sigma}_{\text{Seif}}} \frac{1}{2\pi i} F_\theta = \langle c_1(w), [\hat{\Sigma}_{\text{Seif}}] \rangle =: k,$$

and we were supposing $k$ to be an odd integer. Here again, $\hat{\Sigma}_{\text{Seif}}$ denotes a capped off Seifert surface. The connection $\theta$ has been chosen such that the
curvature $F_\theta$ has compact support inside $S^1 \times D^2$, hence

$$
\int_{\Sigma_{\text{Seif}}} \frac{1}{2} F_\theta = \int_{\{u\} \times D^2} \frac{1}{2} F_\theta
$$

for any $u \in S^1$. On the other hand, if we denote by $\beta$ the integral on the right hand side, then the holonomy $\text{Hol}_l(A)$ around a longitude $l = \{u\} \times \partial D^2 \subseteq \partial(S^1 \times D^2)$ of $K$ is therefore equal to

$$
\text{Hol}_l(A) = \begin{bmatrix}
  e^{i\beta} & 0 \\
  0 & e^{-i\beta}
\end{bmatrix} = \begin{bmatrix}
  e^{i\pi k} & 0 \\
  0 & e^{-i\pi k}
\end{bmatrix}
$$

up to conjugation. Here the first equality is a standard computation of the holonomy of an abelian connection, and the second equation follows from the computations just before. Thus the holonomy of $A$ around the two boundary curves lies on the line $C \subseteq \mathbb{R}(T^2)$. The claim now follows from Proposition 2.1 with the notation taken from Theorem 4.2. 

\[\square\]

6. Donaldson’s invariants for instanton moduli spaces with large holonomy perturbations

The proof of our main theorem will rely essentially on a non-vanishing theorem of Kronheimer-Mrowka [31] about Donaldson’s polynomial invariants of a symplectic 4-manifold $X$ which admits the 0-surgery $Y_0(K)$ of a non-trivial knot $K$ as a separating hypersurface. We start by recalling Donaldson’s invariant derived from moduli spaces of instantons. We then use a variant of the instanton equations which use holonomy perturbations on a neck $[-L, L] \times Y_0(K)$ which are compatible with the holonomy perturbations that we have studied on $Y_0(K)$ before. We show that Donaldson’s invariant can be computed with these large perturbations, essentially by showing that, under certain assumptions on the perturbed representation varieties over $Y_0(K)$, the moduli spaces of instantons contain no reducibles over 1-parameter family of holonomy perturbations, provided one chooses $L$ large enough. Finally, we derive a vanishing result of Donaldson’s invariants of $X$ for possible counter examples of our main theorem.

6.1. Review of instanton gauge theory and Donaldson’s invariants. For the material in this subsection the books of Donaldson and Kronheimer [17], and of Freed and Uhlenbeck [23] can be taken as general references. We assume we have given a Hermitian rank-2 bundle $E_X \to X$ on a Riemannian
4-manifold $X$, with determinant line bundle $v \to X$. We assume a $U(1)$-connection $\theta_X$ in $v$ is fixed. We consider the space $\mathcal{A}_X(E)$ of Hermitian connections on $E_X$ which induce the fixed connection $\theta_X$ on $v$. This is an affine space modelled on $\Omega^1(X; \mathfrak{su}(E))$. To avoid confusion with connections on bundles on 3-manifolds, we write connections on 4-manifolds with bold face letters, so $\mathbf{A}$ denotes a connection in $\mathcal{A}_X$, and $A$ denotes a connection on a 3-manifold.

We denote the group of determinant-1 automorphisms of the bundle $E_X \to X$ by $\mathcal{G}_X$. This group $\mathcal{G}_X$ acts on $\mathcal{A}_X$ in an obvious way on the left. A connection $A$ is called reducible if the stabiliser $\Gamma_A \subseteq \mathcal{G}_X$ of $A$ is strictly bigger than $\mathbb{Z} = \pm \text{id}$, and irreducible otherwise. If there is an $\mathbf{A}$-parallel decomposition of $E$ into two line bundles, and irreducible otherwise.

The Riemannian metric defines the Hodge-star operator $\ast: \Lambda^i(T^*X) \to \Lambda^{4-i}(T^*X)$ on the cotangent bundle, inducing a decomposition into eigenspaces $\Lambda^2(T^*X) = \Lambda^2_+(T^*X) \oplus \Lambda^2_-(T^*X)$ associated to the two eigenvalues $\pm 1$. For a 2-form we denote by $\omega_\pm$ the corresponding orthogonal projection.

The anti-selfduality equation for a connection $\mathbf{A} \in \mathcal{A}_X$ is

$$F_\mathbf{A}^+ = 0.$$  \hspace{1cm} (39)

**Definition 6.1.** The solutions of (39) are called instantons, and the space

$$M_\mathbf{A}^v(X) = \{ [\mathbf{A}] \in \mathcal{A}_X / \mathcal{G}_X \mid F_\mathbf{A}^+ = 0 \}$$

is called moduli space of instantons.

This moduli space of instantons depends on the Riemannian metric and contains reducible connections in general. We denote by $M_\mathbf{A}^v(X)^* \subseteq M_\mathbf{A}^v(X)$ the subspace of instantons which are irreducible. If $b_2^+(X) > 0$, one can achieve that $M_\mathbf{A}^v(X)$ does not contain reducibles for generic metrics. If $b_2^+(X) > 1$, one can achieve that any two metrics with only irreducible instantons can be joined by a path of metrics along which all moduli spaces have only irreducible instantons. Furthermore, the moduli spaces $M_\mathbf{A}^v(X)^*$ are cut out transversally for generic metrics, and in this generic case $M_\mathbf{A}^v(X)^*$ has the structure of a smooth manifold of the expected dimension

$$d = -2\langle p_1(\mathfrak{su}(E)), [X] \rangle + 3(b_1(X) - b_2^+(X) - 1),$$  \hspace{1cm} (40)

where $p_1(\mathfrak{su}(E)) \in H^4(X; \mathbb{Z})$ denotes the first Pontryagin class of the $SO(3)$-bundle $\mathfrak{su}(E)$.  

The moduli spaces $M^v_E(X)$ are in general not compact, but there is a natural compactification due to Uhlenbeck [23] [17].

Associated to 2-dimensional homology classes $a \in H_2(X; \mathbb{Z})$ one can construct codimension-2 submanifolds $V(a)$ of $\mathcal{M}^*_X/G_X$ and these intersect $M^v_E(X)$ in codimension 2 submanifolds generically. To a 0-dimensional homology class $x \in H_0(X; \mathbb{Z})$ one can likewise construct codimension-4 submanifolds $V(x)$. Finally, we mention that the moduli spaces $M^v_E(X)$ are naturally oriented if one fixes an orientation on the real vector space $H^1(X; \mathbb{R}) \oplus H^2(X; \mathbb{R})$. With this at hand, the polynomial invariant of Donaldson’s [15] is a map $D^v: \mathbb{A}(X) \to \mathbb{Q}$, where $\mathbb{A}(X)$ is the symmetric graded algebra generated by $H_0(X; \mathbb{Z}) \oplus H_2(X; \mathbb{Z})$, graded such the grading of $a \in H_1(X; \mathbb{Z})$ is $4 - i$. It is defined for 4-manifolds with $b_2^+(X) > 1$ as follows. If an element $a_1 \otimes a_k \otimes x_1 \otimes x_m \in \mathbb{A}(X)$ has grading $d$ which is the non-negative expected dimension in equation (40) of $M^v_E(X)$ for some bundle $E_X \to X$, then

$$D^v(a_1 \otimes a_k \otimes x_1 \otimes x_m) = \# M^v_E(X)^* \cap \mathcal{V}(a_1) \cap \cdots \cap \mathcal{V}(x_m) \quad (41)$$

for a generic intersection resulting in a compact 0-dimensional manifold, and where the count is taken with the natural orientation. For other elements of $\mathbb{A}(X)$, $D^v$ is taken to be equal to zero.

6.2. Kronheimer-Mrowka’s non-vanishing theorem. In [31], Kronheimer and Mrowka have established a non-vanishing result for quite general symplectic 4-manifolds which contain 3-manifolds admitting a taut foliation as a separating hypersurface, building on work of Feehan-Leness [20], and Eliashberg [18]. By work of Gabai [24] and Eliashberg-Thurston [19] any 0-surgery on a non-trivial knot in $S^3$ admits a taut foliation. We state their theorem only in the special situation that we will need here.

**Theorem 6.2.** Let $K$ be a non-trivial knot in $S^3$. Then the 3-manifold $Y_0(K)$ embeds as a separating hypersurface in a symplectic 4-manifold $X$ with $b_+^v(X) > 1$ for which the following holds:

1. The first homology group $H_1(X; \mathbb{Z})$ is zero.

2. The restriction map $H^2(X; \mathbb{Z}) \to H^2(Y_0(K); \mathbb{Z}) \cong \mathbb{Z}$ is onto.
(3) For any complex line bundle \( v \to X \), Donaldson’s polynomial invariants
\[
D^v_X : \mathcal{A}(X) \to \mathbb{Q}
\]
are non-zero.

6.3. Moduli spaces with holonomy perturbations, neck-stretching arguments. We now study a particular family of metrics on a symplectic 4-manifold \( X \) as in Theorem \ref{thm:6.2}. We start by choosing a Riemannian metric on the 3-manifold \( Y \) (which later will be chosen to be \( Y_0(K) \)), and we choose a Riemannian metric on \( X \) which is a product metric on a tubular neighbourhood of \( Y \). For some \( L > 0 \), we can cut open \( X \) along \( Y \), and insert a cylinder \([-L-1, L+1] \times Y\) with its canonical Riemannian metric into \( X \), resulting in a Riemannian manifold that we denote by \( X(L) \),

\[
X(L) \supseteq [-L - 1, L + 1] \times Y.
\]

We will now choose a smooth cutoff function \( h : X(L) \to [0, 1] \) which is constant 1 on the neck \([-L, L] \times Y\), and which has compact support inside \([-L - 1, L + 1] \times Y\).

We will continue to assume that \( E \to Y \) is a Hermitian bundle on the 3-manifold \( Y \), and we assume that the bundle \( E_X \to X \), when restricted to the neck, is isomorphic to the bundle \( E \to Y \) pulled back to the neck via the projection \([-L - 1, L + 1] \times Y_0(K) \to Y_0(K)\), and we assume we have fixed such an identification in the following. We also assume that the connections in \( \mathscr{A}_X \) are such that the induced connection \( \theta_X \) in the determinant line bundle \( v \) coincides with the pull-back of the fixed connection \( \theta \) in the determinant line bundle \( w \to Y \) on the neck.

Any connection \( A \) on \( E_X \to X \) then yields a family of connection
\[
A(t) := A\big|_{\{t\} \times Y}
\]
on \( E \to Y \) by restriction to the slices \( \{t\} \times Y \). A connection \( A \) is said to be in temporal gauge on the neck if the connection coincides with the trivial product connection if restricted to the lines \( \mathbb{R} \times \{y\} \).

Let \( \{t_k, \chi_k\}_{k=0}^{n-1} \) be holonomy perturbation data on \( Y \) as before. For a connection \( A \in \mathscr{A}_X \) we shall consider the perturbed anti-selfduality equation
\[
F_A^+ = h \cdot \sum_{k=0}^{n-1} \chi'_k(\text{Hol}_{t_k(A)}) \mu_k^+,
\]
(42)
where the self-dual part $\mu^+$ of a 2-form $\mu$ on $Y$ is given by $dt \wedge \ast_3 \mu$, with $\ast_3$ the Hodge-star on the 3-manifold $Y$. The notation is also meant to indicate that at the slice $\{t\} \times Y$ the expression $\text{Hol}_{i_k}(A)$ is meant to be $\text{Hol}_{i_k}(A(t))$.

**Definition 6.3.** The solutions of (42) are called holonomy perturbed instantons, and the space

$$M^+_E(X(L); \{i_k, \chi_k\}) = \{[A] \in \mathcal{A}_X/G_X \mid A \text{ satisfies (42)}\}$$

is called moduli space of perturbed instantons associated to the holonomy perturbation data $\{i_k, \chi_k\}$ and the cutoff function $h$.

We will also need to consider a 1-dimensional family of holonomy perturbations $\theta^s_{i_k, \chi_k}: \mathcal{A} \rightarrow \Omega^2(Y; \text{su}(E))$, parametrised by $s \in [0, 1]$, that interpolates between 0 and the full perturbation $\sum \chi_k' \text{Hol}_{i_k}(A) \mu_k$.

**Definition 6.4.** Given holonomy perturbation data $\{i_k, \chi_k\}^n_{k=0}$ on $Y$ as before, we consider, for $s \in [0, 1]$, the 1-parameter family of perturbations

$$\Phi(s): \mathcal{A} \rightarrow \mathbb{R}$$

of the Chern-Simons function $CS: \mathcal{A} \rightarrow \mathbb{R}$, which for $\frac{k}{n} \leq s \leq \frac{k+1}{n}$ and some $k = 0, \ldots, n-1$ is given by the formula

$$\Phi(s) = n(s - \frac{k}{n}) \Phi_k + \sum_{l=0}^{k-1} \Phi_l.$$

(In particular, $\Phi(0) = 0$, and $\Phi(1) = \sum \Phi_k$.) Here $\Phi_i: \mathcal{A} \rightarrow \mathbb{R}$ denotes the function corresponding to the holonomy perturbation $(i_i, \chi_i)$, as in Section 7 above.

**Proposition 6.5.** The critical points of the perturbed Chern-Simons function

$$CS + \Phi(s): \mathcal{A} \rightarrow \mathbb{R}$$

are the elements $A \in \mathcal{A}$ which solve the equation

$$(FA)_0 = \theta^s_{i_k, \chi_k}(A),$$  \hspace{1cm} (43)$$

where the function $\theta^s: \mathcal{A} \rightarrow \Omega^2(Y; \text{su}(E))$ is defined in the following way. For $\frac{k}{n} \leq s \leq \frac{k+1}{n}$ and some $k = 0, \ldots, n-1$,

$$\theta^s_{i_k, \chi_k}(A) := n(s - \frac{k}{n}) \cdot \chi'_k(\text{Hol}_{i_k}(A)) \mu_k + \sum_{l=0}^{k-1} \chi'_l(\text{Hol}_{i_l}(A)) \mu_l.$$
Proof. This is an immediate consequence of Proposition 1.2. □

Definition 6.6. For choices as above, we denote
\[ R^w(Y; \theta^s_{\{l_k, \chi_k\}}) = \{ [A] \in \mathcal{A}/\mathcal{G} \mid A \text{ solves equation } (43) \}, \]
and we call this the one-parameter perturbed SO(3)-representation variety of \( Y \) associated to \( w \) and the holonomy perturbation data \( \{l_k, \chi_k\} \).

Definition 6.7. If we extend the function \( \theta^s_{\{l_k, \chi_k\}} : \mathcal{A} \to \Omega^2(Y; \mathfrak{su}(E)) \) to the neck as before, we consider the perturbed anti-selfduality equation
\[ F^+_A = h \cdot \theta^s_{\{l_k, \chi_k\}}(A)^+, \quad (44) \]
and we call
\[ M^s_E(X(L); \theta^s_{\{l_k, \chi_k\}}) := \{ [A] \in \mathcal{A}_X/\mathcal{G}_X \mid A \text{ satisfies } (44) \}, \]
the moduli space of instantons associated to the 1-parameter family of holonomy perturbation data \( \theta^s_{\{l_k, \chi_k\}} \), with \( s \in [0, 1] \).

If \( b^+_2(X) > 0 \) the moduli space of instantons \( M^s_E(X) \) contains no reducible connections for a generic Riemannian metric. Furthermore, if \( b^+_2(X) > 1 \) the moduli spaces \( M^s_E(X)(g_s) \) associated to a generic 1-parameter family of Riemannian metrics \( g_s \) contain no reducible connections along the family for a generic 1-parameter family of Riemannian metrics. Presumably, the proof of this fact carries over to the case where we consider only small perturbations along the neck in \( X(L) \). However, this is not the case in our situation of large perturbations, where the standard approach in this situation, see (the proof of) [17, Proposition 4.3.14 and Corollary 4.3.15], clearly cannot be adopted in a straightforward way. We assume that a result of the following kind is implicit in the corresponding situation in [32].

Proposition 6.8 (Avoidance of Reducibles). Suppose that for any \( s \in [0, 1] \) the perturbed representation variety \( R^w(Y; \theta^s_{\{l_k, \chi_k\}}) \) contains no reducible connections. Then the following holds: There is some \( L_0 > 0 \), such that for any \( L \geq L_0 \) and for any \( s \in [0, 1] \) the moduli space
\[ M^s_E(X(L); \theta^s_{\{l_k, \chi_k\}}) \]
contains no reducible connection.
Proof. We follow a strategy similar to the proof of \cite[Proposition 15]{32}, using the 1-parameter family of the perturbed Chern-Simons function in a compactness argument. It uses the relation of the Chern-Simons function to characteristic classes.

Suppose the claim were not true. Then there is a sequence of real numbers \((L_i)_{i \in \mathbb{N}}\) which tends to infinity, some sequence of numbers \((s_i)_{i \in \mathbb{N}}\) in \([0, 1]\) such that the moduli space

\[ M^v_E(X(L_i), \theta_i^{s_i}_{\omega_{\epsilon}}, \xi) \]

contains an equivalence class of a reducible connection represented by \(A_i\).

The expression

\[ e(A_i) = \int_{X_{L_i}} \text{tr}((F_{A_i})_0 \wedge (F_{A_i})_0) = \|F_{A_i}^{-1}\|^2 - \|F_{A_i}^+\|^2 \]

represents the first Pontryagin class of the bundle \(\mathfrak{su}(E)\), and is independent of \(i\). As in \cite{32}, we write \(e(A_i|X')\) for a codimension-0 submanifold \(X'\) of \(X\).

Next we realise that the holonomy perturbation expression on the right hand side of equation (44) is uniformly \(L^\infty\)-bounded by some constant \(C\) which can be taken independent of \(s\) (it can be chosen equal to the \(L^\infty\)-norm of the 2-form \(\mu\) involved in the construction of the holonomy perturbation).

Hence one has a uniform lower bound

\[ e(A_i|[-L_i - 1, -L_i] \cup [L_i, L_i + 1]) \geq -C^2 \text{vol}(Y)^2. \]

As the holonomy perturbations have support in the cylinder \([-L_i - 1, L_i + 1] \times Y\), one also has the lower bound

\[ e(A_i|X(L_i) \setminus [-L_i - 1, L_i + 1] \times Y) \geq 0 \]

for all \(i\), as \((F_{A_i}^+)_0 = 0\) over this piece. Hence, as in \cite{32} Proposition 15] one obtains an uniform upper bound on the neck,

\[ e(A_i|[-L_i - 1, L_i + 1] \times Y) \leq K \]

for some constant \(K\) which is independent of \(s\).

A fundamental fact of instanton Floer theory is that the equations (44) for the perturbed instantons \(A_i\), restricted to the neck \([-L_i, L_i] \times Y\), put in temporal gauge, take the shape of a downward flow equation

\[ \frac{dA_i(t)}{dt} = -\text{grad}(\text{CS} + \Phi(s_i))(A_i(t)), \quad (45) \]

where here again \(A_i(t)\) denotes the restriction of \(A_i\) to the slice \(\{t\} \times Y\). Therefore, \(\text{CS} + \Phi(s_i)\) is monotone decreasing along the path \(t \to A_i(t)\).
On the other hand, it is a well-known fact that the difference of the Chern-Simons function is related to the relative Pontryagin class, and so we have

$$\text{CS}(A_i(-L_i)) - \text{CS}(A_i(L_i)) = e(A_i[-L_i, L_i] \times Y) \leq K.$$ 

The function $\Phi(s)$ has a uniform bound

$$|\Phi(s)| \leq K'$$

for all $s \in [0, 1]$ on $\mathcal{A}$ which only depends on the class functions $\chi_k$ involved in the holonomy perturbation data $\{\iota_k, \chi_k\}$. Hence one obtains a uniform bound on the total drop on the neck as in loc. cit.

$$(\text{CS} + \Phi(s_i))(A_i(-L_i)) - (\text{CS} + \Phi(s_i))(A_i(L_i)) \leq K + 2K'.$$

Given any $\delta > 0$, there is a sequence of intervals

$$(a_i, b_i) \subseteq [-L_i, L_i]$$

of length $\delta$ so that the drop of $\text{CS} + \Phi(s_i)$ from $\{a_i\} \times Y$ to $\{b_i\} \times Y$ converges to 0 as $i$ goes to infinity.

Up to passing to a subsequence, we may suppose that the sequence $(s_i)$ converges to a limit $s_0$. By making $\delta$ small enough so that the $L^2$-norm of the curvature is small enough on the smaller necks $(a_i, b_i) \times Y$ (this can be done uniformly and independent of $s$), one can apply Uhlenbeck’s compactness theorem (by working, after a translation, on the fixed strip $(0, \delta) \times Y$, first on balls, but then by the usual patching process on the entire strip), with the following conclusion. There is a subsequence such that the sequence of connections $A_i$ converges in $C^\infty$ to a limit connection $A$ on $(0, \delta) \times Y$. By continuity, this limit connection is a critical point of the Chern-Simons function $\text{CS} + \Phi(s_0)$. Hence the constant path $(0, \delta) \to \mathcal{A}$ given by $t \mapsto A(t) = A|_{(t) \times Y}$ represents an element in the perturbed representation variety,

$$[A(t)] \in R^w(Y; \theta^{s_0}_{\{\iota_k, \chi_k\}}).$$

By assumption, this is an irreducible connection. However, this implies that for $i$ large enough the nearby connections $A_i(t)$ also are irreducible, as irreducibility is an open condition (even in the weaker $L^\infty$-norm on $\mathcal{A}$.) Hence we have obtained a contradiction. \qed
Remark 6.9 (Uhlenbeck compactification). The Uhlenbeck compactification of moduli spaces carries over to moduli spaces with holonomy perturbations of the kind considered here. This means that the moduli spaces with holonomy perturbations admit natural Uhlenbeck-compactifications analogously to the classical situation. For a discussion, we refer to Donaldson's book on instanton Floer homology [14, Section 5.5]. Therefore, invariants can be defined as in the standard situation, provided one can deal with the reducibles. In particular, one can define the number in (41) with the moduli space $M^v_E(X(L); \theta^s_{\{\iota_k, \chi_k\}})$ instead of $M^v_E(X(L))$.

Proposition 6.10 (Vanishing result). Suppose that the perturbed representation variety $R^w_{\{\iota_k, \chi_k\}}(Y)$ is empty. Then there is some $L_0 > 0$, such that for any $L \geq L_0$ the moduli space $M^v_E(X(L); \{\iota_k, \chi_k\})$ is empty.

Let $b^+_2(X) > 1$. If we assume in addition that the 1-parameter family of holonomy perturbation data $\theta^s_{\{\iota_k, \chi_k\}}$ is such that for any $s \in [0, 1]$ the perturbed representation variety $R^w(Y; \theta^s_{\{\iota_k, \chi_k\}})$ contains no reducible connections, then Donaldson’s polynomial invariant $D^v: \mathbb{A}(X) \to \mathbb{Q}$, is constant to zero.

Proof. The proof of the first part works completely analogously to the proof of Proposition 6.8.

For the second part, the assumption that for all $s \in [0, 1]$ there are no reducibles in $R^w(Y; \theta^s_{\{\iota_k, \chi_k\}})$ implies that for large enough $L$ the family of moduli spaces $M^v(X(L); \theta^s_{\{\iota_k, \chi_k\}})$ contain no reducibles for all $s \in [0, 1]$ by Proposition 6.8. The moduli space $M^v(X(L); \theta^s_{\{\iota_k, \chi_k\}})$ for $s = 0$ corresponds to a moduli space without holonomy perturbations. With such a moduli space, Donaldson’s polynomial invariant is defined, and it is independent of the choice of Riemannian metric (to achieve transversality, perturbations of the metric in some ball away from the neck are enough.) \qed

Remark 6.11. The assumption $b^+_2(X) > 1$ is not necessary in the previous Proposition. It just makes the statement simpler, and sufficient for what we need. Otherwise a similar statement holds for Donaldson’s invariant of a metric in a particular chamber.
7. The main theorem

Our main result is the following

**Theorem 7.1.** Let \( K \) be a non-trivial knot in \( S^3 \). Then the image of \( R(K) \) in the cut-open pillowcase \( C = [0, \pi] \times (\mathbb{R}/2\pi\mathbb{Z}) \) contains an embedded curve which is homologically non-trivial in \( H_1(C; \mathbb{Z}) \cong \mathbb{Z} \).

As stated in the introduction, this is a consequence of the following theorem

**Theorem 7.2.** Let \( K \) be a non-trivial knot. Then any embedded path from \( P = (0, \pi) \) to \( Q = (\pi, \pi) \) in the pillowcase, missing the line \( \{ \beta = 0 \pmod{2\pi} \mathbb{Z} \} \), has an intersection point with the image of \( R(K) \).

**Proof of Theorem 7.1, assuming Theorem 7.2.** The image of \( R(K) \) in the pillowcase \( R(T^2) \) is a compact semi-algebraic variety of dimension \( \leq 1 \). By a theorem of Whitney’s, this is an embedded finite graph \( \Gamma \), see [7] for details about real algebraic geometry. By the properties of the image \( \Gamma \) of \( R(K) \) in \( R(T^2) \) listed in the introduction, the points \( P \) and \( Q \) lie in the complement of the compact subspace graph \( \Gamma \). Theorem 7.2 states that both points \( P \) and \( Q \) lie in different path components of the complement of \( \Gamma \) in the pillowcase \( R(T^2) \). The statement now follows from the following Lemma if we cap off the cut-open pillowcase \( C = [0, \pi] \times (\mathbb{R}/2\pi\mathbb{Z}) \) with two disks on the two boundary curves (or by reglueing the cut-open pillowcase.)

**Lemma 7.3.** Let \( \Gamma \subseteq S^2 \) be a finite embedded graph in the 2-sphere, and let \( P \) and \( Q \) be two points in its complement. Suppose that \( P \) and \( Q \) do not lie in the same path component of \( S^2 \setminus \Gamma \). Then \( \Gamma \) contains an embedded closed curve \( \gamma \) which is homologically non-trivial in \( S^2 \setminus \{ P, Q \} \simeq S^1 \).

**Proof.** We may assume that \( \Gamma \) is connected and still separates \( P \) and \( Q \). We consider a surface with boundary \( F \subseteq S^2 \) which we define to be the set of points which have distance \( \leq \varepsilon \) for some small \( \varepsilon \), chosen so that \( F \) deformation-retracts onto \( \Gamma \), and so that \( P \) and \( Q \) still lie in the complement of \( F \). Its boundary is a disjoint union of circles in \( S^2 \) (which may have ‘corners’.)

Any of these boundary circles bounds a disk in \( S^2 \setminus F^o \). In fact, if there were a connected component \( \Sigma \) of \( S^2 \setminus F^o \) different from a disk, it would have at least two boundary components. This contradicts the assumed connectedness of \( \Gamma \) by the Jordan curve theorem.
We construct a new surface with boundary $\overline{F}$ which we obtain by adding every disk of $S^2 \setminus F^\circ$ to $F$ which does not contain $P$ or $Q$. By Alexander duality,
\[ \tilde{H}_0(S^2 \setminus \overline{F}) \cong H_1(F) \cong \mathbb{Z}, \]
where $\tilde{H}$ denotes reduced singular homology. Therefore, $\overline{F}$ contains a homologically non-trivial embedded closed curve $\gamma$ representing a generator of $H_1(F)$. Without loss of generality, we may assume that $\gamma$ lies inside $F$, because we have only added disks to obtain $\overline{F}$, and we may go further and assume that $\gamma$ lies inside $\Gamma$.

We claim that $\gamma$ is homologically essential in $S^2 \setminus \{P, Q\}$. In fact, $\gamma$ decomposes $S^2$ into two (closed) disks with boundary $\gamma$ by the Jordan curve theorem. Suppose that one of the two disks contained both $P$ and $Q$. Then the other disk must lie entirely in $\overline{F}$, by construction. This contradicts the fact that $\gamma$ is a generator of $H_1(\overline{F})$. Therefore, $P$ and $Q$ lie on different sides of $\gamma$. This clearly implies that $\gamma$ is homologically essential in $S^2 \setminus \{P, Q\}$. □

This terminates the Proof of Theorem 7.1, assuming Theorem 7.2. □

We denote by $\omega$ the standard area form on the 2-torus $T = \mathbb{R}^2/2\pi \mathbb{Z}^2$ (with Euclidean structure) given by $\omega = dx \wedge dy$, where $x, y$ are the coordinates of $\mathbb{R}^2$ in the standard basis. The area form $\omega$ is also a symplectic form, and an area-preserving diffeomorphism of $T$ is the same as a symplectomorphism of the symplectic manifold $T$. We also notice that $\omega$ is invariant under the hyperelliptic involution $(x, y) \mapsto (-x, -y)$ of the torus.

**Definition 7.4.** When we have a family $(\phi^{(t)})_{t \in I}$ with $I \subseteq \mathbb{R}$ of smooth self-maps of a smooth manifold $M$, then we shall say it is a smooth family if the induced map
\[ \phi: I \times T \to T \]
\[ (t, p) \mapsto \phi^{(t)}(p) \]
is smooth of class $C^\infty$.

The author is thankful to Thomas Vogel for providing a sketch of proof of the following Lemma which uses the ‘Moser-trick’.
Lemma 7.5. Let $(\varphi(t))_{t \in [0,1]}$ be a smooth isotopy of the 2-torus $T$ such that $\varphi(0) = id$, and let $c \subseteq T$ be a homologically essential simple closed curve. Then there is a smooth family of symplectomorphisms $(\psi(t))_{t \in [0,1]}$ such that $\psi(0) = id$, and such that we have

$$\varphi(t)(c) = \psi(t)(c)$$

for all $t \in [0,1]$. (The last equality is an equality of sets, $\varphi(t)$ and $\psi(t)$ do not have to coincide point-wise on $c$.) If the family $(\varphi(t))$ is $\mathbb{Z}/2$-equivariant, then we can arrange that the family $(\psi(t))$ is $\mathbb{Z}/2$-equivariant as well.

Proof. Up to scaling $\varphi(t)$ in directions normal to $c$, we may suppose that

$$\left(\varphi(t)\right)^*\omega_{\mid c} = \omega_{\mid c}$$

for all $t \in [0,1]$. For any $t$, we define a 1-parameter family of area forms on $T$,

$$\omega_s^{(t)} := s\omega + (1 - s)(\varphi(t))^*\omega,$$

for $s \in [0,1]$. We notice that this family is constant along $c$ in the parameter $s$ because of assumption (46). Furthermore, the cohomology class $[\omega_s^{(t)}]$ in the 2nd de Rham cohomology group is also constant in $s$. Therefore, the derivative with respect to $s$ is an exact 2-form. Hence there is a family of 1-forms $(\alpha_s^{(t)})_{s \in [0,1]}$ such that we have

$$\frac{d\omega_s^{(t)}}{ds} = -d\alpha_s^{(t)}$$

for all $s \in [0,1]$. Furthermore, we may suppose that $\alpha_s^{(t)}$ depends smoothly on $s$ and $t$. In fact, the Hodge decomposition theorem states that one has an $L^2$-orthogonal decomposition

$$\Omega^2(T) = \Delta\Omega^2(T) \oplus \mathcal{H}^2(T),$$

where $\Delta = dd^* : \Omega^2(T) \to \Omega^2(T)$ is the Laplace-Beltrami operator, and where $\mathcal{H}^2(T)$ denotes the (one-dimensional) space of harmonic 2-forms. Therefore, for an exact 2-form $\beta$ we have $\beta = dd^* \beta$, and the 1-form $\alpha = d^* \beta$ depends smoothly on any parameter $\beta$ depends on, and satisfies $d\alpha = \beta$.

Because the 2-forms $\omega_s^{(t)}$ are non-degenerate, there is a unique smooth 1-parameter family of vector fields $X_s^{(t)}$ on $T$ satisfying

$$\omega_s^{(t)}(X_s^{(t)}, \cdot) = \alpha_s^{(t)}$$

for all $s \in [0,1]$. We claim that we may, and from now on will suppose that the vector field $X_s^{(t)}$ is tangent to $c$ for all $s$ and all $t$. 
In fact, by choosing our coordinates appropriately, we may suppose that the simple closed curve $c$ is defined, say, by the equation $y = \pi$ in our coordinates $(x, y)$ on $T$, 

$$c = \{(x, y) \mid y = \pi\}.$$ 

We may write the form $\alpha^{(t)}_s$ as 

$$\alpha^{(t)}_s = a^{(t)}_s(x, y) \, dx + b^{(t)}_s(x, y) \, dy.$$ 

Then the form 

$$\tilde{\alpha}^{(t)}_s := \alpha^{(t)}_s - a^{(t)}_s(x, \pi) \, dx$$ 

depends also smoothly on $s$ and $t$, its differential is equal to $-\frac{d\omega^{(t)}_s}{ds}$ as well, and $\tilde{\alpha}^{(t)}_s$ has vanishing $dx$-component along $c$. But this means that the vector field $X^{(t)}_s$ has vanishing $\frac{\partial}{\partial y}$-component for all $s$ and $t$ along $c$.

As $T$ is closed, the $s$-dependent vector field $X^{(t)}_s$ may be integrated to an isotopy $(\phi^{(t)}_s)$ for all time $s$. More precisely, we have

$$\frac{d\psi^{(t)}_s}{ds}((\phi^{(t)}_s)^{-1}(p)) = X^{(t)}_s(p)$$

for all point $p \in T$, and for all $s$ and $t$. It is standard to verify that $\phi^{(t)}_s$ depends smoothly on $t$, if $\alpha^{(t)}_s$ does.

We claim that the family of 2-forms $(\phi^{(t)}_s)^*\omega^{(t)}_s$ is constant in $s$. In fact, if we denote by $\mathcal{L}$ the Lie derivative along a vector field, we have \[1\]

$$\frac{d}{ds} (\phi^{(t)}_s)^*\omega^{(t)}_s = (\phi^{(t)}_s)^* (\mathcal{L}_{X^{(t)}_s} \omega^{(t)}_s) + \frac{d\omega^{(t)}_s}{ds}$$

$$= (\phi^{(t)}_s)^* (i_{X^{(t)}_s} d\omega^{(t)}_s + d i_{X^{(t)}_s} \omega^{(t)}_s) + \frac{d\omega^{(t)}_s}{ds}$$

$$= (\phi^{(t)}_s)^* (d\alpha^{(t)}_s + \frac{d\omega^{(t)}_s}{ds})$$

$$= 0$$

for all $s$ and $t$, by Cartan’s formula $\mathcal{L}_{X} \beta = i_X (d\beta) + d(i_X \beta)$. In particular, we have

$$(\phi^{(t)}_1)^* \omega^{(t)}_1 = (\phi^{(t)}_0)^* \omega^{(t)}_0 = \omega^{(t)}_0 = (\varphi^{(t)})^* \omega$$

for all $t$. But as $\omega^{(t)}_1 = \omega$ for all $t$, this gives

$$(\phi^{(t)}_1)^* \omega = (\varphi^{(t)})^* \omega$$
and hence 
\((\varphi^{(t)} \circ (\phi_1^{(t)})^{-1})^* \omega = \omega\)

for all \(t\).

The smooth family of symplectomorphisms
\[ \psi^{(t)} : = \varphi^{(t)} \circ (\phi_1^{(t)})^{-1} \]
now has the desired property. In fact, as the vector field \(X_s^{(t)}\) restricted to \(c\) is tangent to \(c\) for all \(s\) and \(t\), the map \(\phi_1^{(t)}\) maps \(c\) to itself for all \(t\).

Finally, if \(\varphi^{(t)}\) is \(\mathbb{Z}/2\)-equivariant for all \(t\), (then in particular \(c\) is \(\mathbb{Z}/2\)-invariant,) then the whole argument carries through in a \(\mathbb{Z}/2\)-equivariant way, because the symplectic form \(\omega\) and also the interpolations \(\omega_s^{(t)}\) defined in the proof are \(\mathbb{Z}/2\)-invariant, the operators \(d\) and \(d^*\) anti-commute with the involution when applied to \(\mathbb{Z}/2\)-invariant functions or 2-forms (this is because \(\tau^* dx = -dx\) and \(\tau^* dy = -dy\), hence the 1-forms \(\alpha_s^{(t)}\) and \(\tilde{\alpha}_s^{(t)}\) are \(\mathbb{Z}/2\)-equivariant. Therefore, the vector fields \(X_s^{(t)}\) are equivariant, and hence the integration \(\phi_s^{(t)}\) is is \(\mathbb{Z}/2\)-equivariant. \(\square\)

**Remark 7.6.** The assumption that the curve \(c\) is homologically essential is essential. For homologically inessential curves, the area enclosed by the curves \(c\) is an obvious invariant under area preserving isotopies, and hence the corresponding statement cannot hold in general, as a smooth isotopy \(\varphi^{(t)}\) may change the area enclosed in the curves \(\varphi^{(t)}(c)\). The assumption is slightly hidden in the proof. It is used in the statement that we may suppose that the curve \(c\) is defined by the equation \(y = \pi\).

We now ready to prove Theorem 7.2.

**Proof of Theorem 7.2.** Suppose this were not the case. Then there is an embedded path \(c_1\) from \(P = (\pi, 0)\) to \(Q = (\pi, \pi)\) in \(R(T^2)\) which has empty intersection with \(R(K)\) (a path as the blue one in Figure 2 above). This is smoothly isotopic to the straight line \(c_0 : [0, 1] \rightarrow R(T^2)\), given by \(t \mapsto (\pi, t \cdot \pi)\) through isotopies which keep the four singular points fixed, and we can assume that the isotopy misses the straight line \(d\) from \((0, 0)\) to \((\pi, 0)\) in \(R(T^2)\) for all times.

We lift the problem to the branched double cover \(\tilde{R}(T^2)\). We obtain a \(\mathbb{Z}/2\)-equivariant isotopy \(\varphi_t : \tilde{R}(T^2) \rightarrow \tilde{R}(T^2)\) from \(\varphi_0 = \text{id}\) to \(\varphi_1\), where \(\varphi_1\) maps the lift \(\check{c}_0\) of the straight line \(c_0\) to the lift \(\check{c}_1\) of the \(\mathbb{Z}/2\)-invariant closed curve which is the lift of \(c_1\). Furthermore, for all \(t \in [0, 1]\) the image
of $\varphi_t(\hat{c})$ misses the lift of the curve $d$. By Lemma 7.5 there is a smooth $\mathbb{Z}/2$-equivariant isotopy $\psi_t$ through area-preserving maps for which we have $\varphi_t(\hat{c}_0) = \psi_t(\hat{c}_0)$ for all $t \in [0, 1]$.

By assumption the lift $\hat{R}(K)$ does not intersect $\hat{c}_1$. The subset $\hat{R}(K) \subseteq \hat{R}(T^2)$ is compact. Hence there is some $\varepsilon > 0$ such that the $\varepsilon$-neighbourhood of the image of $\hat{c}_1$ also has empty intersection with $\hat{R}(K)$. We may suppose that this $\varepsilon$ is chosen small enough so that for all $t \in [0, 1]$ the image $\psi_t(\hat{c}_0)$ has distance from the image of $\hat{d}$ at least $\varepsilon$.

By our technical main result Theorem 4.2 we can find a $\mathbb{Z}/2$-equivariant isotopy $\varphi_t$ such that $\varphi_t$ is $\varepsilon$-close to $\psi_t$ for all $t \in [0, 1]$, and which we can realise through holonomy perturbations: There is some holonomy perturbation data $\{\iota_k, \chi_k\}$ on the trivial $SU(2)$-bundle over the thickened torus $M = [0, 1] \times S^1 \times S^1$ such that the restrictions $r_\pm$ to the two boundary components

$$r_\pm: R_{\{t_k, \chi_k\}}(M) \to R(T^2)$$

satisfy $r_+ = \overline{\varphi}_1 \circ r_-$, where $\overline{\varphi}_1$ is the map induced by $\varphi_1$. The holonomy perturbation data $\{t_k, \chi_k\}$ over the thickened torus $M$ also determines holonomy perturbation data on the 0-surgery $Y_0(K)$ of $K$, as in Section 5 above.

By Proposition 5.3 the fact that

$$\varphi_1(\hat{c}_0) \cap \hat{R}(K) = \emptyset$$

implies that the perturbed representation variety is empty,

$$R^w_{\{t_k, \chi_k\}}(Y_0(K)) = R(K|\overline{\varphi}_1(\hat{c}_0)) = \emptyset.$$

This, however, contradicts Kronheimer-Mrowka’s non-vanishing Theorem 6.2 by the stretching argument in Proposition 6.10, because for all $t \in [0, 1]$ the curve $\varphi_t(\hat{c}_0)$ misses the line $\hat{d}$, and hence $R^w(Y_0(K); \theta^t_{\{t_k, \chi_k\}})$ does not contain reducibles for all $t \in [0, 1]$.

8. $SU(2)$-REPRESENTATIONS OF SPlicings OF KNOT COMPLEMENTS

For a knot $K$ in $S^3$ and an open tubular neighbourhood $n(K)$ we denote by $Y(K)$ the complement $S^3 \setminus n(K)$.
We will need one more property about the image of $R(K)$ in the pillowcase manifested as the representation variety of the boundary torus, with the choice of coordinates as beforehand.

Proposition 8.1. (i) If $\rho(t)$ is a path of irreducibles in $R(K)^*$ which has an end point in a reducible connection $\rho_0$, then the image of $\rho_0$ in the pillowcase can be represented by a point $(\alpha, 0)$ where $\Delta_K(e^{2i\alpha}) = 0$, in other words, $e^{i2\alpha}$ is a root of the Alexander polynomial of the knot $K$. As for any knot $K$ one has $\Delta_K(1) = \pm 1$, $\rho_0$ cannot map to a point represented by $(0, 0)$ or $(\pi, 0)$.

(ii) The lines \{\(\alpha = 0 \mod 2\pi\mathbb{Z}\}\} and \{\(\alpha = \pi \mod 2\pi\mathbb{Z}\}\} only contain the two central representations, both with \{\(\beta = 0 \mod 2\pi\mathbb{Z}\}\}. There is a neighbourhood of these two lines which do not contain images of irreducible connections.

Proof. The first property is due to Klassen [29, Theorem 19]. The two lines \{\(\alpha = 0 \mod 2\pi\mathbb{Z}\}\} and \{\(\alpha = \pi \mod 2\pi\mathbb{Z}\}\} cannot contain any non-central representations, as the meridian normally generates the knot group. The second property now follows from the first and the compactness of $R(K)$.

Definition 8.2. Let $K, K'$ be two knots in $S^3$, and let $m_K, l_K$ be meridian and longitude of $K$, and likewise $m_{K'}, l_{K'}$ be meridian and longitude of $K'$. These curves $m_K, l_K$ form a basis of $H^1(T; \mathbb{Z})$ of the boundary torus $T = \partial Y(K)$, and analogously for the boundary torus $T' = \partial Y(K')$. Let $\varphi : T \rightarrow T'$ be homeomorphism such that in the above bases of $H_1(T)$ respectively $H_1(T')$, one has

$$\varphi_* = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$  

Then the glued up 3-manifold $Y_{K,K'} := Y(K) \cup_\varphi Y(K')$ is an integer homology 3-sphere that we call the (untwisted) splicing of $K$ and $K'$.

If instead we have

$$\varphi_* = \begin{pmatrix} 0 & 1 \\ 1 & k \end{pmatrix} \quad \text{or} \quad \varphi_* = \begin{pmatrix} k & 1 \\ 1 & 0 \end{pmatrix}$$

for some non-zero integer $k \in \mathbb{Z}$, we call the 3-manifold $Y_{K,K'} := Y(K) \cup_\varphi Y(K')$ a twisted splicing of $K$ and $K'$.

Our main result, Theorem 7.1 yields the following
Theorem 8.3. Let $K, K'$ be two non-trivial knots in $S^3$.

(i) Then there is an irreducible representation $\rho : \pi_1(Y_{K,K'}) \to SU(2)$ from the fundamental group $\pi_1(Y_{K,K'})$ of the untwisted splicing into $SU(2)$. It restricts to irreducible representations of the two knot groups $\pi_1(Y(K))$ and $\pi_1(Y(K'))$.

(ii) Any twisted splicing $Y_{K,K'}$ of $K$ and $K'$ has an irreducible representation from $\pi_1(Y_{K,K'})$ into $SU(2)$.

Proof. We only prove the first statement, as this is the only one we need later, leaving the second one as an exercise to the reader.

With our conventions from above, Theorem 7.1 implies that the image of the $SU(2)$-representation variety of $R(K)$ in the representation variety of the boundary torus $R(\partial Y(K)) = R(T^2) = (\mathbb{R}^2/2\pi\mathbb{Z}^2)/\tau$, the pillowcase, contains an embedded curve $\gamma_K$ which is homologically non-trivial in the cut-open pillowcase $[0, \pi] \times \mathbb{R}/2\pi\mathbb{Z}$. We consider a lift to the torus $\mathbb{R}^2/2\pi\mathbb{Z}^2$. In the fundamental domain $[0, 2\pi] \times [0, 2\pi]$ a lift of the curve $\gamma_K$ can be represented by a path from the line $\{0\} \times [0, 2\pi]$ to the line $\{2\pi\} \times [0, 2\pi]$. Because of the second property of Proposition 8.1, there is some $\delta(K) > 0$ such that $\gamma_K$ can be represented by a curve which lies in $[\delta(K), \pi - \delta(K)] \times [0, 2\pi]$. A corresponding statement holds for a curve $\gamma_{K'}$ in the pillowcase of the other boundary torus $R(\partial Y(K'))$.

By the Seifert-van-Kampen theorem, representations $\rho_K \in R(K)$ and $\rho_{K'} \in R(K')$ extend to a representation $\rho : \pi_1(Y_{K,K'}) \to SU(2)$, if they coincide on the boundary torus along which they are glued together with $\varphi$. We have the following restriction maps:

$$R(\partial Y(K)) \to R(\partial Y(K)) \xleftarrow{\varphi^*} R(\partial Y(K')) \leftarrow R(K')$$

In our identification of $R(\partial Y(K))$ and $R(\partial Y(K'))$ with $R(T^2)$, the first coordinate of $R(T^2)$ corresponds to the meridian, and the second to the longitude. With these identifications, $\varphi^*$ just swaps the two coordinates,

$$\varphi^*: (\alpha, \beta) \mapsto (\beta, \alpha).$$

Hence, the curve $\varphi^*\gamma_{K'}$ is represented by an embedded path from the line $[0, 2\pi] \times \{0\}$ to the line $[0, 2\pi] \times \{2\pi\}$.

The two paths representing $\gamma_K$ and $\varphi^*\gamma_{K'}$ therefore necessarily intersect in a point of the pillowcase $R(\partial Y(K))$ corresponding to irreducible representations of $R(K)$ and $R(K')$, as an intersection point must have coordinates in one of the squares $[\delta(K), \pi - \delta(K)] \times [\delta(K'), \pi - \delta(K')]$. □
9. $SL(2, \mathbb{C})$-REPRESENTATIONS OF HOMOLOGY 3-SPHERES

The author has learned the following result and its proof from Michel Boileau. It builds on strong results of Boileau, Rubinstein and Wang [8] about the domination of 3-manifolds by maps of degree 1. This was a major motivation for establishing Theorem 7.1 because of the application Theorem 9.4 below.

**Theorem 9.1.** Let $Y'$ be an integer homology sphere different from the 3-sphere. Then there exists a map $Y' \to Y$ of degree 1 onto an integer homology sphere $Y$ which is either hyperbolic, Seifert-fibred or the untwisted splicing of two non-trivial knots in $S^3$.

**Proof.** Suppose

$$Y' = Y^{(0)} \to Y^{(1)} \to Y^{(2)} \to \ldots$$

is an – a priori infinite – sequence of integer homology 3-spheres which are different from the 3-sphere, where the 3-manifolds $Y^{(i)}$ and $Y^{(j)}$ are not homeomorphic to each other if $i \neq j$, and where each map is of degree 1. By Corollary 1 of [8], this sequence must be finite. Let $Y$ be a 3-manifold where the sequence stops. Therefore, if there is a map of degree 1 from $Y$ to another 3-manifold, then the other 3-manifold must be either $Y$ itself or $S^3$. Suppose $Y$ is neither hyperbolic nor Seifert-fibred. Then by the Geometrization Theorem for 3-manifolds [30, 39, 40, 41] there must be an incompressible torus $T$ sitting inside $Y$, splitting $Y$ into 3-manifolds with boundary $Y_1$ and $Y_2$,

$$Y = Y_1 \cup_T Y_2.$$  

Both $Y_1$ and $Y_2$ must be integer homology solid tori, and either injection map $H_1(T) \to H_1(Y_i)$, $i = 1, 2$, must be surjective and have kernel of rank 1. Let $l_i$ be a simple closed curves on $T$ which generates the kernel of this injection map, $i = 1, 2$. It is easy to see that $l_1$ and $l_2$ must have intersection number ±1, hence they form a basis of $H_1(T)$.

The pinching construction of [8, Section 5] yields a degree 1 map of $Y_1$ to an actual solid torus $S^1 \times D^2$ that can be extended by the identity on $Y_2$ to a map of degree 1

$$f : Y \to S^1 \times D^2 \cup_T Y_2 = \overline{Y},$$

where the solid torus is glued in so that the boundary of $D^2$ is homologous to $l_1$. 
We claim that the resulting 3-manifold $\overline{Y}$ cannot be homeomorphic to $Y$. If $\overline{Y}$ were homeomorphic to $Y$, we would have an epimorphism $f_* : \pi_1(Y) \to \pi_1(Y)$.

However, fundamental groups of 3-manifolds are Hopfian (this is because 3-manifold groups are residually finite, see for instance [4]), meaning that any self-epimorphism is an isomorphism, so that $f_*$ would necessarily be an isomorphism. The fundamental group $\mathbb{Z}^2$ of the torus $T \subseteq Y$ injects into $\pi_1(Y)$, because it is incompressible. However, by the pinching construction of the map $f$, the image $f_*(\pi_1(T))$ is not isomorphic to $\mathbb{Z}^2$ since the torus $f(T)$ bounds a solid torus in $\overline{Y}$, hence we obtain a contradiction. Therefore, $\overline{Y}$ must be the 3-sphere.

This means that $Y_2$ is the complement of a tubular neighbourhood of a non-trivial knot $K_2$ in $S^3$. Arguing symmetrically, we see that $Y_1$ is the complement of a tubular neighbourhood of a non-trivial knot $K_1$ in $S^3$.

The curves $l_1$ is a longitude to the knot $K_1$, and the curve $l_2$ is so for $K_2$. As glueing in the solid torus on either side resulted in the 3-sphere, we see that $l_1$ is glued to a meridian of $K_2$, and $l_2$ is glued to a meridian of $K_1$. Hence $Y$ is an untwisted splicing of the knots $K_1$ and $K_2$. □

The $SU(2)$-representation varieties of Seifert fibred homology spheres has been described by Fintushel and Stern in terms of linkages, see [21]. Presumably from this description, one can prove the following, which certainly has been known to experts.

**Proposition 9.2.** Let $Y$ be a Seifert fibered integer homology 3-sphere. Then there is an irreducible representation $\rho : \pi_1(Y) \to SU(2)$.

**Proof.** We give a short proof which is independent of Fintushel and Stern’s description. By Montesinos’ results, every Seifert fibered integer homology 3-sphere $Y$ is the branched double cover $\Sigma_2(K)$ of a Montesinos’ knot $K$ of determinant 1, see [37]. By Kronheimer and Mrowka’s non-vanishing result [33, Corollary 7.17], there is an irreducible representation $\theta : \pi_1(S^3 \setminus K) \to SU(2)$ which maps a meridian of the knot $K$ to an element of trace zero in $SU(2)$. By rather elementary means (see for instance [47, Proposition 9.1]), this implies that $Y = \Sigma_2(K)$ has an irreducible representation of its fundamental group in $SU(2)$. □
Proposition 9.3. Let $Y$ be a hyperbolic integer homology 3-sphere. Then it admits an irreducible representation $\rho: \pi_1(Y) \to SL(2, \mathbb{C})$.

Proof. This follows immediately from the definition, and the fact that the group $PSL(2, \mathbb{C}) = SL(2, \mathbb{C})/\{\pm id\}$ is the orientation preserving isometry group of hyperbolic 3-space. Because $Y$ has no 2-dimensional cohomology with $\mathbb{Z}/2$-coefficients, there is no obstruction to lifting the metric representation to $SL(2, \mathbb{C})$. \qed

We are now ready to prove our main application.

Theorem 9.4. Let $Y$ be an integer homology 3-sphere different from the 3-sphere. Then there is an irreducible representation $\rho: \pi_1(Y) \to SL(2, \mathbb{C})$.

Proof. By Theorem 9.1 above, the three cases considered in Proposition 9.2, Proposition 9.3 and Theorem 8.3 imply the general case, since a degree one map induces an epimorphism of fundamental groups. \qed

10. The higher dimensional case

As stated in the introduction, we have the following result, contrasting Theorem 9.4.

Theorem 10.1. In any dimension $m \geq 4$ there is an integer homology sphere $W^m$ different from the $m$-sphere $S^m$ such that the fundamental group $\pi_1(W)$ has only the trivial representation in $SL(2, \mathbb{C})$.

Proof. Kervaire’s main theorem in [28] states that any perfect group $\pi$ which has trivial second homology group $H_2(\pi)$ (such groups are also called superperfect) is the fundamental group of an integer homology sphere $W^m$ of dimension $n \geq 5$. Hence we are left with the problem of finding such a group $\pi$ which does not admit irreducible representations in $SL(2, \mathbb{C})$. By the following Proposition, the Schur cover of any finite non-abelian simple group except of $A_5$ is such an example. In fact, non-abelian simple groups are perfect, and their Schur covers are superperfect.

The case of dimension 4 requires a more detailed look at Kervaire’s argument. Suppose a presentation of a superperfect group $\pi$ is given to us. One first constructs an $m$-dimensional manifold with boundary $V$ which has one 0-handle, a 1-handle for each generator of the presentation, and a 2-handle
for each relation. The Euler characteristic \( \chi(V) \) is equal to the deficiency of the presentation plus 1. The doubling of \( V \) is the \( m \)-manifold 
\( Z := V \cup \overline{V} \), where \( \overline{V} \) carries the opposite orientation of \( V \), and where the glueing map along the boundary is taken to be the identity. Then \( Z \) has fundamental group \( \pi \), see [28]. In the higher-dimensional case one can perform surgery on \( Z \) in order to obtain \( W \), where \( W \) has no more non-trivial homology except in dimension 0 and \( m \).

If \( m = 4 \), however, \( Z \) is already an integer homology sphere if the deficiency of the chosen presentation of \( \pi \) is 0. In fact \( \chi(V) = 1 \) implies that \( \chi(Z) = 2 \). As \( b_1(Z) = 0 \), we must therefore have \( b_2(Z) = 0 \) also. Furthermore, there is no torsion in the second homology because of the universal coefficient theorem, and because \( H_1(Z; \mathbb{Z}) = 0 \).

The claim for \( m = 4 \) now follows also from the following Proposition because there are many Schur covers of finite non-abelian simple groups which are known to have a presentation of deficiency 0, see for instance the table in [11]. Examples in this table include the infinite families \( SL(2, \mathbb{Z}/p) \) for \( p \) an odd prime, as well as the Schur covers of the groups \( A_6, A_7, Sz(8) \), or the superperfect groups \( PSL(3, 3), PSU(3, 3), M_{11} \). \( \square \)

**Proposition 10.2.** The Schur covers of all finite non-abelian simple groups with the exception of the alternating group \( A_5 \) (the rotation preserving symmetry group of the icosahedron) have only trivial representations in \( SL(2, \mathbb{C}) \).

**Proof.** Let \( G \) be a finite simple group, and let \( C \) be its Schur cover [44]. (Schur covers of perfect groups are unique, and it may be that \( C = G \).) We then have a short exact sequence

\[
1 \to H_2(G; \mathbb{Z}) \to C \to G \to 1,
\]

where \( H_2(G; \mathbb{Z}) \) is central in \( C \).

Suppose we are given an irreducible representation \( \rho: C \to SL(2, \mathbb{C}) \). This has to map \( H_2(A_n) \) to the centre \( \{ \pm \text{id} \} \) of \( SL(2, \mathbb{C}) \). Hence it induces a **projective** representation

\[
\overline{\rho}: G \to PSL(2, \mathbb{C}) = SL(2, \mathbb{C})/\{ \pm \text{id} \}.
\]

The image \( \overline{\rho}(G) \) is a finite subgroup, and hence it must lie in the subgroup \( SO(3) = PSU(2) \subseteq PSL(2, \mathbb{C}) \). As \( G \) is simple, the homomorphism \( \overline{\rho} \) maps \( G \) to either the trivial subgroup of \( SO(3) \), or it maps \( G \) isomorphically to a subgroup. The latter case cannot happen since the only simple non-abelian finite subgroup of \( SO(3) \) is \( A_5 \), which we have excluded by assumption.
We conclude that $\bar{\rho}$ must be the trivial representation. Hence $\rho$ must be a central representation, and hence the trivial one, as the Schur cover $C$ is a perfect group.

□

11. 3-sphere recognition is in coNP, modulo the generalized Riemann hypothesis

Rubinstein has established an algorithm recognizing the 3-sphere from a triangulation of a given 3-manifold [42] which was subsequently simplified by Thompson [46]. Schleimer has shown that the 3-sphere recognition problem lies in the complexity class $\text{NP}$, if one takes as input data a triangulation description of the 3-manifold to start with.

The complexity class $\text{NP}$ consists of the problems that have polynomial time algorithms on non-deterministic Turing machines. Equivalently, there exists a verifier for the problem who can decide whether a proposed solution to the problem indeed is a solution within polynomial time in terms of the size of the input data.

Kuperberg [34] has shown that the assertion that a knot $K \subseteq S^3$ is not the unknot lies in the complexity class $\text{NP}$, provided the generalized Riemann hypothesis (GRH) holds. His result builds on Kronheimer and Mrowka's results in [32] and Theorem [34, Theorem 3.4] that he establishes:

**Theorem 11.1** (Kuperberg). Let $G$ be an affine algebraic group over $\mathbb{Z}$ and assume that GRH holds. Then there is a polynomial $P$ with the following significance: Let $\Gamma$ be a discrete group with a finite presentation of length $l$. If there is a homomorphism

$$\rho_C : \Gamma \to G(\mathbb{C})$$

with non-commutative image, then there is also a homomorphism

$$\rho_p : \Gamma \to G(\mathbb{Z}/p)$$

with non-commutative image, for a prime $p$ such that $\log(p) = P(l)$.

We think of our integer homology sphere $Y$ as being given by a Heegaard diagram. From this we can read off a presentation of the fundamental group. If $g$ is the genus of the Heegaard diagram, and if $k$ is the number of intersections in the Heegaard diagram (counted absolutely, and not up to sign), we obtain a presentation of the fundamental group $\pi_1(Y)$ of length $g + k$. 
Theorem 11.2. Let $Y$ be an integer homology 3-sphere, described by a Heegaard diagram. Then the assertion that $Y$ is not the 3-sphere lies in the complexity class $\text{NP}$, provided GRH holds.

Proof. As noted above, we obtain a presentation of the fundamental group $\pi_1(Y)$ of length $l$ which depends polynomially on the size of the input data in terms of the Heegaard diagram. By Theorem 9.4, there exists a representation $\rho_C : \pi_1(Y) \to SL(2, \mathbb{C})$ with non-abelian image. If GRH holds, then by Theorem 11.1 there exists a prime number $p$ such that $\log(p)$ depends polynomially on the length $l$ of the presentation, for some universal polynomial, and there is a homomorphism

$$\rho_p : \pi_1(Y) \to SL(2, \mathbb{Z}/p)$$

with non-abelian image. A certificate (or proof or witness) for the problem consists of such a prime $p$ and such a representation $\rho_p$.

The verifier checks that $\rho_p$ indeed is a homomorphism by checking that the relations hold. The effort for this is again polynomial in terms of the input data. The verifier then checks that the commutator of at least two different generators is non-trivial. The effort for this is again polynomial in terms of the input data. \hfill \Box

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