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Induced Magnetic Field in sp-Hybridized Carbon Rings: Analysis of Double Aromaticity and Antiaromaticity in Cyclo[2n]carbon Allotropes

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Abstract

The induced magnetic field of C_{2N} (N = 3-14) carbon rings was dissected to contributions from out-of-plane and in-plane π orbitals revealing two concurrent long range shielding or deshielding cones as a manifestation of the dual aromatic and antiaromatic character of C_{4n+2} and of C_{4n} rings respectively. Aromaticity based on the magnetic criterion was evaluated with regard to bonding pattern and geometrical characteristics elucidating the influence of bond length and bond angle alternation on out-of-plane and in-plane magnetic response. Ground state polyynic geometries of C_{4n+2} rings exhibit comparable shielding cones to annulenes, decreasing magnetic response with regard to ring size and similar π_{out} and π_{in} diatropicity. Transition state cumulenic rings display increased aromaticity expressed by a constant very strong magnetic response and augmented π_{out} diatropicity with regard to π_{in}. The variations of the induced magnetic field are explained on the basis of frontier orbital interactions through rotational excitations, which enable further rationalization of the aromatic/antiaromatic behavior.
Introduction

Since the discovery of buckminsterfullerene C_{60} in 1985, research on \( sp^2 \) carbon allotropes such as fullerenes, graphenes and carbon nanotubes have sparked the design and synthesis of a plethora of novel carbon-based nanomaterials with exceptional diversity and applications in numerous fields such as optoelectronics, gas separation, energy conversion and shortage, catalysis and medicine. Around the same time research on different carbon allotropes with mixed \( sp^2-sp \) hybridization or solely \( sp \) hybridized carbon atoms has paved the way for discovery of new all carbon species with promising applications such as graphdiynes and carbon rings.

Early in 1989 Diederich et al. reported the gas phase generation of highly reactive C_{18} cyclocarbon in which each carbon atom is bonded to two other carbons, in contrast to \( sp^2 \) hybridized allotropes where the bonding pattern involves three-coordinated carbon atoms. The main question that evolved to a controversy by conflicting theoretical studies was the bonding pattern of C_{18}: is it polyynic with altering triple and single bonds, or cumulenic with equal bond lengths? Pure and hybrid DFT as well as MP2 calculations predict cumulenic structures, whereas HF, quantum Monte Carlo and CCSD calculations predict polyynic geometries. Further studies diagnosed that the DFT calculated bonding pattern depends on the amount of HF exchange of the hybrid functional, whereas range-separated exchange nonempirical DFT schemes provide polyynic structures regardless the type of functional that is used. Finally in 2019 Kaiser et al. isolated by on-surface synthesis and structurally characterized C_{18} using atomic force microscopy (AFM), revealing a polyynic structure with altering single-triple CC bonds setting an end to the controversy. Moreover cumulenic structure has been predicted to be a transition state between two degenerate polyynic ground states with inverted single-triple bonds.

A unique characteristic of \( sp \)-hybridized carbon structures is the presence of two perpendicular \( \pi \)-electron systems extending the capabilities from \( sp^2 \)-hybridized carbon structures which sustain a single \( \pi \)-system. Diederich proposed that C_{18} is benefited from two distinctive sets of delocalized \( \pi \) electrons, one from out-of-plane (\( \pi_{\text{out}} \)) and the other from in-plane (\( \pi_{\text{in}} \)) oriented 18\( \pi \) orbitals (Figure 1), both obeying Hückel’s 4n+2 rule, which is expected to exhibit interesting electron acceptor properties. The presence of double aromaticity for C_{4n+2} and antiaromaticity for C_{4n} cyclocarbons was confirmed via ring current analysis by Fowler et al., while the large NICS
values reported by Suresh\textsuperscript{20} also imply the existence of double aromaticity given by the two perpendicular $\pi$-systems. The double aromaticity of small carbon rings and related boron-carbon molecules was also confirmed with CMO-NICS analysis by Schleyer \textit{et al.}\textsuperscript{32} However studies on the aromaticity of carbon rings were performed on cumulenic structures and the effect of bonding pattern on the magnetic response was not evaluated. Recently Baryshnikov \textit{et al.}\textsuperscript{28} employed the GIMIC method to measure the ring currents of both polyynic and cumulenic configurations of C\textsubscript{18} and reported that the polyynic structure induces a strong diatropic current which is however much weaker than the diatropic current of cumulene. Moreover, although GIMIC does not allow yet the dissection of ring current to contributions from different sets of orbitals, they obtained the current strengths of $\pi\text{\textsubscript{out}}$ and $\pi\text{\textsubscript{in}}$ sets by integrating the current density within distinctive contours and found that $\pi\text{\textsubscript{out}}$ orbitals induce stronger diatropic currents than $\pi\text{\textsubscript{in}}$, both in polyynic and cumulenic geometry.

\textbf{Figure 1.} Illustration of two orthogonal sets of $p$ orbitals of C\textsubscript{18} ring displaying perpendicular to the plane ($\pi\text{\textsubscript{out}}$) and in-plane ($\pi\text{\textsubscript{in}}$) orientation.

Herein, we study the distinctive characteristics of double aromaticity and antiaromaticity based on the magnetic criterion\textsuperscript{33} from cyclo[2N]carbon allotropes, given by their peculiar perpendicular $\pi$-systems. However we note that a thorough investigation of aromaticity requires to take into consideration electronic and energetic criteria as well.\textsuperscript{34–36} We perform a detailed analysis of the magnetic response of C\textsubscript{2N} (N = 3-14) rings directly related to aromatic/antiaromatic behavior,\textsuperscript{37–42} by means of canonical orbitals contributions to the induced magnetic field (CMO-IMF) which allows the through-space visualization of shielding and deshielding cones induced by different sets of orbitals, as well as from each CMO.\textsuperscript{43–46} Hence the magnetic field induced from each $\pi\text{\textsubscript{out}}$ and $\pi\text{\textsubscript{in}}$ sets was visualized for polyynic and cumulenic geometries of C\textsubscript{4n+2} rings, as well as for C\textsubscript{4n}
rings, revealing the trends of the magnetic response of each set of MOs and the effect of ring size, bond length and bond angle. Our results show that polyynic geometries induce a cumulative magnetic field from both sets which is comparable to annulenes and declines with the evolution of ring size, whereas cumulenes induce a much stronger magnetic field independent to the ring size. Accordingly, $C_{4n}$ rings display decreasing deshielding cones for both sets with regard to ring size. These trends are associated to the inherent bonding pattern and geometrical characteristics of carbon rings, which are further explained on the basis of frontier orbital interactions through rotational excitations.

**Computational Methods**

A number of various exchange–correlation functionals (XC) were tested for the optimization of $C_{18}$ (See Supporting Information, Table S1). Functionals with high percentage of HF exchange ($\text{HFX}>50\%$) yielded polyynic structures, whereas pure GGAs and functionals with low percentage of HFX resulted in cumulenic structures with no bond length alteration (BLA) (Table S1). In this work we have chosen to use the range-separated hybrid $\omega$B97XD for the optimization of the $C_{2N}$ ($N = 3-14$) molecules because it correctly reproduces the experimentally realized polyynic geometry of $C_{18}$ and it has been reported to have the best performance among other functionals for the description of polyyinic structures. Moreover, long-range corrected XC functionals like $\omega$B97XD have been found to provide a better description of cyclic electron delocalization in monocyclic and polycyclic aromatics with regard to conventional XC functionals. All geometry optimizations were carried out with Gaussian09 using the 6-311++G(d, p) gaussian basis set.

Nucleus independent chemical shieldings (NICS) were calculated within the GIAO formalism using several functionals, as detailed in Supporting Information, and the triple-$\zeta$ slater type basis set with two polarizations functions (TZ2P) employing the ADF2019 software. BH\text{andHLYP} was chosen for the quantitative analysis of dissected NICS$_{zz}$ because it performs in very good agreement with $\omega$B97XD (Table S2, Figure S2) and it is available for CMO analysis of the chemical shift in ADF. The computational affordable PBE was chosen for the qualitative analysis of the induced magnetic field. Although PBE was found to overestimate NICS, especially in polyyinic structures, it accurately reproduces the qualitative features of the magnetic response with regard to the ring size and bonding pattern of $C_{2N}$ rings (see Supporting Information, Figures S2-
S3). Chemical shieldings were dissected to contributions from canonical MOs (CMOs) employing the NMR and EPR modules\textsuperscript{54} of ADF and the contributions of $\pi_{\text{out}}$ and $\pi_{\text{in}}$ sets of orbitals were constructed by the summation of the corresponding CMOs.

For the calculation of the induced magnetic field the molecules were placed on xy Cartesian plane and the chemical shielding was calculated in a $31^3$ cubic grid of points with a separation step of 0.5Å. Extraction of CMO contributions at each point and generation of cube files per MO was performed with our custom MIMAF code.\textsuperscript{55} Visualization of induced magnetic field isosurfaces was performed with VMD.\textsuperscript{56}

**Results and Discussion**

**Geometrical characteristics**

The ground state optimized geometries of $C_{4n+2}$ $(n=1-6)$ and $C_{4n}$ $(n=2-6)$ molecules are shown in Figure S1 and their corresponding geometrical parameters are given in Table 1. All molecules adopt $C_{(n/2)h}$ point group symmetry representing polyynic structures with altering bond lengths and altering bond angles, except for $C_6$ and $C_{10}$ which adopt $D_{(n/2)h}$ cumulenic structures with equal bond lengths and altering angles. The bond length alteration (BLA) ranges from 0.123Å to 0.148Å for all $C_{(n/2)h}$ point group molecules except for $C_{14}$ which presents a minimal BLA of 0.051Å. On the other hand, bond angle alteration (BAA) is very large for small rings starting from 58.4º and 55.6º for $C_6$ and $C_8$ respectively, and decreases rapidly as the ring size increases, reaching to a constant value of 1.7º for $C_{2N}$ $(N=11-14)$, while the smallest BAA (0.4º) is observed in $C_{18}$. For $C_{4n}$ molecules $D_{(n/2)h}$ geometries with altering bond lengths and equal bond angles were found to be first-order saddle points, whereas higher symmetry $D_{nh}$ geometries with equal bonds and angles were found to be higher-order saddle points for both $C_{4n+2}$ and $C_{4n}$ molecules.
Table 1. Geometrical parameters of $C_{2n}$ (N = 3-14) rings optimized at the ωB97XD/6-311++G(d, p) level.

| Molecule | Point Group | $d_1$ (Å) | $d_2$ (Å) | BLA$^a$ | $\theta_1$ (°) | $\theta_2$ (°) | BAA$^b$ |
|----------|-------------|-----------|-----------|---------|---------------|---------------|---------|
| $C_6$    | $D_{3h}$    | 1.321     | 0.00      | 149.2   | 90.8          | 58.4          |
| $C_8$    | $C_{4h}$    | 1.384     | 1.251     | 0.133   | 162.8         | 107.2         | 55.6    |
| $C_{10}$ | $D_{5h}$    | 1.290     | 1.290     | 0.00    | 165.0         | 123.0         | 42.1    |
| $C_{12}$ | $C_{6h}$    | 1.359     | 1.234     | 0.126   | 168.3         | 131.7         | 36.6    |
| $C_{14}$ | $C_{7h}$    | 1.307     | 1.256     | 0.051   | 167.0         | 141.6         | 25.4    |
| $C_{14}$ TS | $D_{3h}$    | 1.281     | 0.00      | 168.1   | 140.5         | 27.6          |
| $C_{16}$ | $C_{8h}$    | 1.364     | 1.216     | 0.148   | 158.9         | 156.1         | 2.8     |
| $C_{18}$ | $C_{9h}$    | 1.346     | 1.223     | 0.123   | 160.2         | 159.8         | 0.4     |
| $C_{18}$ TS | $D_{3h}$    | 1.277     | 0.00      | 169.4   | 150.6         | 18.8          |
| $C_{20}$ | $C_{10h}$   | 1.358     | 1.216     | 0.141   | 162.6         | 161.4         | 1.2     |
| $C_{22}$ | $C_{11h}$   | 1.351     | 1.219     | 0.132   | 164.5         | 162.8         | 1.7     |
| $C_{22}$ TS | $D_{11h}$   | 1.276     | 0.00      | 171.0   | 156.3         | 14.8          |
| $C_{24}$ | $C_{12h}$   | 1.355     | 1.216     | 0.139   | 165.8         | 164.2         | 1.7     |
| $C_{26}$ | $C_{13h}$   | 1.352     | 1.217     | 0.135   | 167.0         | 165.3         | 1.7     |
| $C_{26}$ TS | $D_{13h}$   | 1.276     | 0.00      | 172.7   | 159.6         | 13.1          |
| $C_{28}$ | $C_{14h}$   | 1.354     | 1.216     | 0.138   | 168.0         | 166.3         | 1.7     |

$^a$Bond length alteration, $d_1$-$d_2$. $^b$Bond angle alteration, $\theta_1$-$\theta_2$.

For $C_{4n+2}$ (n=3-6) molecules, QST2 and IRC calculations generated cumulenic geometries with equal bond lengths and altering bond angles, adopting $D_{(n/2)h}$ point group symmetry, as transition states between two inversion related polyynic ground states$^{28}$ (Figure 2). The transition states display increased BAA with regard to the respective ground states. The energy barrier for the inversion increases with the number of carbon atoms, starting from a marginal barrier of 0.12 kcal·mol$^{-1}$ for $C_{14}$ and reaching to 10.90, 24.42 and 37.54 kcal·mol$^{-1}$ for $C_{18}$, $C_{22}$ and $C_{26}$ respectively.
Figure 2. Relative IRC profiles of C\textsubscript{18}, C\textsubscript{22} and C\textsubscript{26} depicting the interchange of single-triple bonds of polyynic ground states through a cumulenic transition state.

Total magnetic response

When an (anti)aromatic molecule is subjected into a uniform external magnetic field $B^\text{ext}$, then closed circuits of mobile electrons induce a current density, which in turn induce a secondary magnetic field, $B^{\text{ind}}$, either opposing (shielding) the external field in aromatic rings, or reinforcing (deshielding) the external field in antiaromatic rings. In planar structures the molecule is oriented such as the molecular plane (xy) is perpendicular to the external field’s direction (z), and hence the induced magnetic field is designated as $B^{\text{ind}}_z$. The induced magnetic field is typically visualized as isosurfaces of $B^{\text{ind}}_z$, forming long range shielding or deshielding cones in aromatic or antiaromatic rings respectively. The value of $B^{\text{ind}}_z$ isosurface at any point is equal to the zz tensor component of nucleus independent chemical shift, NICS\textsubscript{zz}.

NICS\textsubscript{zz} values, presented in Table 2 and Figure S2, reveal the characteristic features of carbon rings’ magnetic response. These features are: a) C\textsubscript{4n+2} molecules display large diatropic (negative) NICS\textsubscript{zz} representative of aromatic character, whereas C\textsubscript{4n} molecules display large paratropic (positive) NICS\textsubscript{zz} representative of antiaromatic character; b) NICS\textsubscript{zz} of polyynic C\textsubscript{4n+2} (n=3-6) and C\textsubscript{4n} molecules decreases significantly as the ring size increases; c) NICS\textsubscript{zz} of the cumulenic C\textsubscript{4n+2} (n=1-6) geometries increases with the ring size, showing a trend to reach a maximum value at larger rings; d) accordingly for C\textsubscript{4n+2} molecules there is an increasing difference of NICS\textsubscript{zz}. 

between polyynic ground states and cumulenic transition states with the increment of ring size; e) the maximum NICS\(_{zz}\) among ground state C\(_{4n+2}\) molecules is observed in C\(_{14}\) which presents the minimum BLA.

The above remarks clarify controverting NICS studies of C\(_{4n+2}\) rings.\(^{19,20}\) NICS values reported by Fowler et al. show an increment of diatropic NICS of C\(_{4n+2}\) for n=1-5 and then a significant decrement for n=6,7 due to molecular geometry obtained with B3LYP functional which predicts cumulenic structures for n=1-5 and polyynic structures for n=6,7.\(^{19}\) On the other hand, Suresh and Remya reported increasing diatropic NICS values for all C\(_{4n+2}\) (n=1-7) trending to a maximum value due to cumulenic geometries obtained with MO6L functional.\(^{20}\)

**Table 2.** Total NICS\(_{zz}\) values (ppm) and contributions from \(\pi_{\text{out}}\) and \(\pi_{\text{in}}\) orbitals of C\(_{2N}\) carbon rings and reference molecules computed at the BHandHLYP/TZ2P level.

| Molecule | Total | \(\pi_{\text{out}}\) | \(\pi_{\text{in}}\) | Molecule | Total | \(\pi_{\text{out}}\) | \(\pi_{\text{in}}\) |
|----------|-------|-----------------|-----------------|----------|-------|-----------------|-----------------|
| C\(_6\)   | -37.7 | -39.8           | 19.2            | C\(_8\)  | 130.9 | 110.5           | 28.2            |
| C\(_{10}\) | -65.2 | -52.1           | -3.3            | C\(_{12}\)| 96.8  | 64.8            | 34.1            |
| C\(_{14}\) | -81.6 | -49.9           | -25.4           | C\(_{16}\)| 57.9  | 27.0            | 31.7            |
| C\(_{14}\)TS| -89.2 | -56.7           | -25.0           | C\(_{20}\)| 33.4  | 15.5            | 18.7            |
| C\(_{18}\) | -44.7 | -20.6           | -16.4           | C\(_{24}\)| 18.7  | 8.6             | 10.8            |
| C\(_{18}\)TS| -101.5| -58.3           | -37.2           | C\(_{28}\)| 10.5  | 4.8             | 6.4             |
| C\(_{22}\) | -22.0 | -10.7           | -7.8            | C\(_8\)\(_6\)| 107.2 | 90.2            |                 |
| C\(_{22}\)TS| -104.0| -58.1           | -40.7           |         |       |                 |                 |
| C\(_{26}\) | -10.9 | -5.4            | -3.4            |         |       |                 |                 |
| C\(_{26}\)TS| -99.3 | -55.8           | -39.0           |         |       |                 |                 |
| C\(_6\)\(_6\)| -15.9 | -36.7           |                 |         |       |                 |                 |
| C\(_{18}\)\(_18\)| -41.5 | -48.0           |                 |         |       |                 |                 |
| C\(_{18}\)\(_6\)| -8.1  | -12.7           | -15.7           |         |       |                 |                 |
| B\(_9\)\(_9\)| 3.3   | 0.5             | 3.4             |         |       |                 |                 |

Visualizations of B\(_z^{\text{ind}}\) induced by all electrons of C\(_{2N}\) molecules are presented in Figure 3 using multiple clipped isosurfaces (±5ppm and ±20ppm) to depict the long-range effect, as well as a full isosurface of large B\(_z^{\text{ind}}\) values to illustrate specific characteristics of the magnetic response close to the molecular domain. C\(_{4n+2}\) molecules induce strong long-range shielding cones representative of strong aromatic character. The long-range shielding cone (cyan, -5ppm) of C\(_{4n+2}\) molecules
increases with the ring size but the actual trend is revealed by inspecting the strong response close to the molecular plane depicted with green isosurfaces (-50 ppm) in Figure 3. In ground state polyynic structures the extension of the short-range response increases from C\textsubscript{10} to C\textsubscript{14} where it reaches a maximum span forming a uniform shielding cone, and then gradually decreases from C\textsubscript{18} to C\textsubscript{26}, deforming to a toroidal shape inside the ring of C\textsubscript{26}. On the contrary the strong shielding cone of cumulenic transition states retains its uniform long ranged shape. Therefore, the weakening of the magnetic response in ground states is attributable to polyynic bonding and is irrelevant to the ring size. However, the discrepancy of the magnetic response between ground and transition states for C\textsubscript{18}, C\textsubscript{22} and C\textsubscript{26} is apparently evolving with the ring size. Thus, according to the magnetic criterion, the aromaticity of polyynic C\textsubscript{4n+2} decreases for n>3 while it is retained for the respective cumulenic transition states.

Figure 3. Isosurfaces of total $B_z^{\text{ind}}$ of transition states $D_{(n/2)h}$ C\textsubscript{4n+2} (n=3-6) (botom row), ground states $C_{(n/2)h}$ C\textsubscript{4n+2} (n=1-6) (middle row) and C\textsubscript{4n} (n=2-6) (top row).

C\textsubscript{4n} molecules sustain strong long-range deshielding cones representative of strong antiaromatic character. However strong short-range response (+60 ppm) decreases with the evolution of ring
size, starting with extended uniform deshielding cones in C\textsubscript{8} and C\textsubscript{12} and downgrading to toroidal shape in C\textsubscript{24}, which implies weakening of antiaromaticity.

**Magnetic response of \(\pi_{\text{out}}\) and \(\pi_{\text{in}}\) orbitals**

The total magnetic response suggests the (anti)aromatic character of C\textsubscript{2N} carbon rings but does not give clear information about the source of the long range (de)shielding cones and the individual role of each (para)diatropic circuit to the double aromatic character, because it is constructed by contributions from all \(\pi_{\text{out}}, \pi_{\text{in}}, \sigma\) and core orbitals. In order to gain insight about the origins and the duality of the (anti)aromatic character of C\textsubscript{2N} carbon rings, we dissected the magnetic response to contributions from \(\pi_{\text{out}}\) and \(\pi_{\text{in}}\) sets of orbitals, which are the exclusive sources of the strong long-range magnetic response.

![Image of isosurfaces of \(\pi_{\text{out}}\) and \(\pi_{\text{in}}\) contributions to \(B_z^{\text{ind}}\) of polyynic ground states of C\textsubscript{4n+2} (n=1-6) rings.](image)

**Figure 4.** Isosurfaces of \(\pi_{\text{out}}\) (top) and \(\pi_{\text{in}}\) (bottom) contributions to \(B_z^{\text{ind}}\) of polyynic ground states of C\textsubscript{4n+2} (n=1-6) rings.

In Figure 4 the respective magnetic response of \(\pi_{\text{out}}\) and \(\pi_{\text{in}}\) sets of orbitals of ground states C\textsubscript{4n+2} molecules are given. It is clear that for C\textsubscript{14} and larger rings, both \(\pi_{\text{out}}\) and \(\pi_{\text{in}}\) orbitals induce long range shielding cones, illustrating the dual source of their diatropicity. Smaller members C\textsubscript{6} and C\textsubscript{10} display none or very weak in-plane magnetic response respectively, due to large bond angle alteration (BAA) which leads to small in-plane orbital overlap. Specifically, the \(\pi_{\text{out}}\) orbitals of C\textsubscript{6} induce a shielding cone very similar to benzene and have comparable \(\pi_{\text{out}}\)-NICSzz values (-39.8 and -36.7 ppm for C\textsubscript{6} and C\textsubscript{6}H\textsubscript{6} respectively). The \(\pi_{\text{in}}\) orbitals of C\textsubscript{6} induce a weak short range diatropic response with a paratropic sphere inside the ring (+19.2 ppm at ring center), which is
representative of localized CC σ bonding.\textsuperscript{43} Therefore, C\textsubscript{6} displays equal diatropicity with benzene originating from the out-of-plane π system, and the large difference of their total NICS\textsubscript{zz} does not arise from delocalized electrons but from localized σ bonds which have larger paratropic contributions in benzene. C\textsubscript{10} presents a strong long-range shielding cone induced from π\textsubscript{out} orbitals with -52.1 ppm contributions at ring center, and a considerable weaker shielding cone induced from π\textsubscript{in} orbitals with only -3.3 ppm at ring center, denoting a marginal in-plane π magnetic response due to large BAA.

C\textsubscript{14} induces long range shielding cones from both π\textsubscript{out} and π\textsubscript{in} orbitals but the out-of-plane response is much stronger than the in-plane, as π\textsubscript{out} forms an extended uniform -40 ppm isosurface, while π\textsubscript{in} forms a condensed shielding toroid inside the ring and the π\textsubscript{out}-NICS\textsubscript{zz} (-49.9 ppm) is the double of π\textsubscript{in}-NICS\textsubscript{zz} (-25.4 ppm). This difference is attributed to the geometry of C\textsubscript{14} which presents minimum BLA that favors out-of-plane delocalization and significant BAA that hinders in-plane delocalization. The corresponding transition state displays almost identical magnetic response, both in terms of shielding cones and NICS\textsubscript{zz} values, due to the small structural differences of the two states.

For the ground states of larger rings C\textsubscript{18}, C\textsubscript{22} and C\textsubscript{26}, π\textsubscript{out} and π\textsubscript{in} display very similar shielding cones. The π\textsubscript{in}-NICS\textsubscript{zz} is only 2-4 ppm less diatropic than π\textsubscript{out}-NICS\textsubscript{zz}, denoting marginally increased out-of-plane magnetic response. The diatropicity decreases with the increment of the ring size for both sets of orbitals. Specifically NICS\textsubscript{zz} of π\textsubscript{out} (π\textsubscript{in}) decrease from -20.6 (-16.4) ppm in C\textsubscript{18} down to -10.7 (-7.8) and -5.4 (-3.4) ppm in C\textsubscript{22} and C\textsubscript{26} respectively.

In order to assess C\textsubscript{18} diatropicity we used as reference three prototypical molecules with 18-π electrons and varying degree of aromaticity, namely 18-annulene, 18-dehydroannulene and the BN analogue\textsuperscript{57} B\textsubscript{9}N\textsubscript{9}, presented in Figure 5. Compared to the π response of 18-annulene (Figure 5.b), the π\textsubscript{out} of C\textsubscript{18} induces a weaker shielding cone and presents less than half NICS\textsubscript{zz}, denoting a considerable weaker out-of-plane magnetic response in C\textsubscript{18}. If we consider both π\textsubscript{out} and π\textsubscript{in} orbitals of C\textsubscript{18}, their cumulative shielding cone (Figure 5.a) is greater than B\textsuperscript{ind}\textsubscript{zz} of 18-annulene, but this is due to the overestimation of PBE which predicts a total (π\textsubscript{out} + π\textsubscript{in})-NICS\textsubscript{zz} value 21 ppm more diatropic than NICS\textsubscript{zz} of C\textsubscript{18}H\textsubscript{18}, while with BHandHLYP it is 11.0 ppm less diatropic. Therefore, C\textsubscript{18} displays similar or slightly weaker diatropicity compared to C\textsubscript{18}H\textsubscript{18}. Furthermore, 18-
dehydroannulene (Figure 5.c) is less aromatic than C_{18} and induces a weaker $\pi_{\text{out}}$ shielding cone ($\pi_{\text{out}}$-NICS$_{zz}$ is -12.7 ppm), as well as a weak short-range paratropic response from 12 $\pi_{\text{in}}$ electrons ($\pi_{\text{in}}$-NICS$_{zz}$ is +15.7 ppm). Finally, B$_9$N$_9$ (Figure 5.d) is a characteristic example with extreme localized electrons, with both 18 $\pi_{\text{out}}$ and 18 $\pi_{\text{in}}$ electrons inducing weak short range diatropic cones strictly positioned on nitrogen atoms. Hence C$_{4n+2}$ ($n>3$) carbon rings, although they exhibit dual source of diatropicity, their magnetic response is comparable to classical $\pi$-aromatic annulenes.

![Figure 5. Isosurfaces of $\pi$ contributions to B$^{\text{ind}}_{zz}$ of a) C$_{9h}$ C$_{18}$, b) 18-annulene, c) 18-dehydroannulene and d) B$_9$N$_9$.](image)

The magnetic response of transition states, presented in Figure 6, display constant strong shielding cones of both $\pi_{\text{out}}$ and $\pi_{\text{in}}$. Specifically, $\pi_{\text{out}}$-NICS$_{zz}$ retains a value around -57 ppm, whereas $\pi_{\text{in}}$-NICS$_{zz}$ amounts to ~ -39 ppm for all transitions states, implying favorable $\pi_{\text{out}}$ delocalization. The difference of $\pi_{\text{out}}$ and $\pi_{\text{in}}$ can be explained by the increased BAA of cumulenic structures which hinders in-plane overlap. Indeed, the 3$^{rd}$-order saddle point $D_{18h}$ geometry of C$_{18}$ with zero BLA and BAA presents almost equal $\pi_{\text{out}}$ and $\pi_{\text{in}}$ NICS$_{zz}$ values (-60.2 and -58.7 ppm respectively). However, $\pi_{\text{in}}$ electrons in cumulenic geometries are still more diatropic than polyynic. Hence BLA is the main factor that constrains both $\pi_{\text{out}}$ and $\pi_{\text{in}}$ magnetic response and BAA affects secondarily only $\pi_{\text{in}}$ magnetic response. Consequently, the cumulenic geometries are very diatropic from both sources of magnetic response.
Concerning the antiaromatic $C_{4n}$ ($n=2-6$) rings, both $\pi_{\text{out}}$ and $\pi_{\text{in}}$ orbitals induce equivalent long range deshielding cones (Figure 7), except $C_8$ which exhibits only $\pi_{\text{out}}$ deshielding cone and a weak short range paratropic response of $\pi_{\text{in}}$ due to large BAA, which does not justify a paratropic in-plane ring current. The $\pi_{\text{in}}$ response of $C_{14}$ is also weak due to large BAA but still forms a deshielding cone denoting weak in-plane paratropic current, whereas for larger rings $\pi_{\text{out}}$ and $\pi_{\text{in}}$ induce equivalent deshielding cones. The $\pi_{\text{out}}$ deshielding cones of $C_8$ and $C_{12}$ are comparable to planar $C_8H_8$ but decline significantly for larger rings, as the $\pi_{\text{in}}$ deshielding cones also do. Hence, $C_{4n}$ ($n=3-6$) rings exhibit dual source of paratropicity representative of antiaromatic character which weakens significantly with the increment of the ring size.
Figure 7. Isosurfaces of $\pi_{\text{out}}$ (top) and $\pi_{\text{in}}$ (bottom) contributions to $B_z^{\text{ind}}$ of $C_{4n}$ ($n=2-6$) rings and planar $C_8H_8$.

**CMO contributions**

In order to elucidate the above remarks, we have to take into consideration the electronic structure of carbon rings and especially the frontier orbitals which principally determine the overall magnetic response. Dissection of the induced magnetic field to CMO contributions reveals that lower energy valence orbitals induce diatropic response, whereas the highest occupied orbitals dictate the overall magnetic response.\textsuperscript{43,44} Indeed in $C_{2N}$ the contributions of $\pi_{\text{out}}$ and $\pi_{\text{in}}$ to NICS$_{zz}$ are linear correlated to the contribution of the corresponding HOMOs ($R^2 = 0.97$, Figure S4). Hence the (para)diatropicity can be analyzed in terms of HOMO contributions, given in Table S4 and Figures S5-S6.

Generally, in antiaromatic molecules the HOMO induce very strong paratropic response that overwhelms the diatropicity of lower orbitals, whereas in aromatic molecules the HOMOs induce very weak paratropic or diatropic response that adjusts the overall diatropicity and tune the aromaticity. In turn, the paratropic response of HOMOs originates from symmetry allowed rotational excitations to unoccupied orbitals and its magnitude depends on the energy gap and the overlap of interacting orbitals.\textsuperscript{43,44,46,58–60}

In $C_{4n}$ rings the HOMO and LUMO of both out-of-plane and in-plane orientation, are non-degenerate MOs with the same number of $n$ nodal planes and the same symmetry. Hence the symmetry allowed HOMO$_{\text{out/in}}$→LUMO$_{\text{out/in}}$ excitations, rotating the HOMOs by an angle of
leads to optimum overlap producing maximum paratropic response. For example, in C\textsubscript{16} the rotational excitations of HOMO\textsubscript{out} and HOMO\textsubscript{in} with 4 nodal planes rotated by 22.5° leads to perfect overlap with LUMO\textsubscript{out} and LUMO\textsubscript{in} contributing with +63.3 and +64.4 ppm respectively to NICS\textsubscript{zz} (Figure 8.b) and induce very strong long range deshielding cones (Figure 8.a) that dominate on the overall magnetic response. As the HOMO-LUMO gaps remain practically unchanged throughout the C\textsubscript{4n} series (Table S4), the decline of paratropicity originates from weakening paratropic response of HOMOs and is attributed to the decrease of rotational overlap as the number of nodal planes increases.

In C\textsubscript{4n+2} rings, both HOMO\textsubscript{out} and HOMO\textsubscript{in} are doubly degenerate with \( n \) nodal planes, whereas the LUMOs are doubly degenerate with \( n+1 \) nodal planes. In a perfect symmetry of \( D_{nh} \) point group a HOMO→LUMO rotational excitation would be symmetry forbidden, but in \( D_{(n/2)h} \) and \( C_{(n/2)h} \) point groups such excitations are symmetry allowed. However due to different number of nodal planes the overlap is small, inducing weak magnetic response that depends on the geometrical characteristics of the ring. For example in \( C_{9h} \) \( C_{18} \), the frontier orbitals are distorted due to significant BLA and the HOMO→LUMO excitations display a small overlap inducing weak paratropic response (Figure 8.a) that diminish the diatropicity of the lower energy orbitals, contributing with +27.8 and +30.6 ppm to NICS\textsubscript{zz} of \( \pi_{\text{out}} \) and \( \pi_{\text{in}} \) respectively (Figure 8.b). In contrast, in \( D_{9h} \) \( C_{18} \) the zero BLA leads to negligible rotational overlap of HOMO\textsubscript{out}→LUMO\textsubscript{out} excitation, resulting to diatropic response of HOMO\textsubscript{out} that adds to the overall diatropicity. On the other hand, HOMO\textsubscript{in}→LUMO\textsubscript{in} excitations of \( D_{9h} \) \( C_{18} \) induce small paratropic contributions (+15.9 ppm) due to increased BAA that lead to weaker diatropicity of \( \pi_{\text{in}} \) with regard to \( \pi_{\text{out}} \). The same holds for the ground states of smaller members \( C_6 \), \( C_{10} \) and \( C_{14} \) with zero (or almost zero) BLA, which induce shielding cones from HOMO\textsubscript{out} (Figure S5) and display increased \( \pi_{\text{out}} \) diatropicity. For C\textsubscript{4n+2} ground states the increment of HOMO nodal planes with the ring size results to the gradual increment of the small overlap and hence to the augmentation of HOMO paratropic contributions (Figure S5), causing the gradual decrement of the overall diatropicity.
Figure 8. a) Isosurfaces of HOMO\textsubscript{out} and HOMO\textsubscript{in} contributions to $B_z^{\text{ind}}$ of C\textsubscript{16}, ground state C\textsubscript{9h} C\textsubscript{18} and transition state D\textsubscript{9h} C\textsubscript{18}. b) Energy levels and contributions of HOMO\textsubscript{out/in}→LUMO\textsubscript{out/in} rotational excitations of C\textsubscript{16}, ground state C\textsubscript{9h} C\textsubscript{18} and transition state D\textsubscript{9h} C\textsubscript{18}.

Conclusions

The dissection of the induced magnetic field to contributions from out-of-plane and in-plane $\pi$ orbitals revealed the dual aromatic and antiaromatic character of C\textsubscript{4n+2} and C\textsubscript{4n} rings respectively, according to the magnetic criterion. The magnetic response induced from distinctive sets of electrons is sensitive to the bonding pattern and geometrical features, where polyynic C\textsubscript{4n+2} ground states display similar out-of-plane and in-plane magnetic response and comparable diatropicity to annulenes, while the cumulenic transition states exhibit very strong diatropic character and augmented out-of-plane magnetic response with regard to in-plane magnetic response. Bond length alteration is the prime factor that constrains both $\pi$\textsubscript{out} and $\pi$\textsubscript{in} magnetic response, whereas bond angle alteration affects secondarily only $\pi$\textsubscript{in} magnetic response. A significant decline of the magnetic response with the increment of ring size is observed for both ground state aromatic and
antiaromatic rings, which is attributed to decreasing paratropic contributions of HOMOs originating from rotational excitations.

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