On the global structure of distant galactic disks

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Abstract. Radial and vertical profiles are determined for a sample of 34 edge-on disk galaxies in the HDFs, selected for their apparent diameter larger than 1.3\arcsec and their unperturbed morphology. The thickness and flatness of their galactic disks are determined and discussed with regard to evolution with redshift. We find that sub-$L^*$ spiral galaxies with $z \sim 1$ have a relative thickness or flatness (characterized by $h_z/h$ the scaleheight to scalelength ratio) globally similar to those in the local Universe. A slight trend is however apparent, with the $h_z/h$ flatness ratio larger by a factor of $\sim 1.5$ in distant galaxies if compared to local samples. In absolute value, the disks are smaller than in present-day galaxies. About half of the $z \sim 1$ spiral disks show a non-exponential surface brightness distribution.

Key words. galaxies: evolution – galaxies: high-redshift – galaxies: structure

1. Introduction

A very fundamental characteristic of galactic disks is the flatness, defined by the ratio between exponential scale-lengths in the vertical and radial direction $h_z/h$. The determination of the thickness of stellar disks is a difficult enterprise and has been determined mostly for edge-on galaxies, although attempts have been made in other cases, too (Ma et al. 1998). Typical values are found to be between 0.15 and 1 (van der Kruit & Searle 1981, 1982; de Grijs 1998) with a systematic trend for galaxies to become thinner from S0 to Sc if a more complex multicomponent structure (thick/thin disk) is not taken into account. Stellar velocity dispersion measurements show that the dispersion tends to grow with total luminosity, and vary radially as the square root of surface density, so as to support a constant scaleheight (Bottema 1993). Recently Kregel et al. (2002) have shown that the flattening of disks increases with the amplitude of their maximum rotation, and their total HI mass.

This flatness is an important clue to the formation and evolution history of galaxies. Disk stars are thought to be formed out of a very thin gaseous disk. The observed structure then thickens due to scattering by local fluctuations of the gravitational field as caused, e.g., by giant molecular clouds or spiral structure in the disk. The velocity dispersion of the stars increases in both the radial and vertical direction by this diffusion through phase space (Wielen 1977, Binney & Lacey 1988). The disk may thicken more efficiently through interactions with small companions or minor mergers, as suggested by numerical simulations (Quinn et al. 1993, Walker et al, 1996, Velazquez & White 1999). The thickening could amount to a factor of 1.5–2 and can be a constraint on the frequency of mergers if disks are observed today to be too thin (e.g. Toth & Ostriker 1992). The existence of a thick stellar disk in our own Galaxy has been attributed to an old merger event (Robin et al 1996, Dalcanton & Bernstein 2002). However, other phenomena have to be taken into account since the thin disk could also be maintained through gas infall with subsequent star formation. Many observational facts strongly suggest a high gas infall rate (e.g. Toomre, 1990, Sancisi et al. 1990, Jiang & Binney 1999). Dynamics of galaxy disks, bar and spiral reformation constrain this rate such that galaxies may double their mass in less than 10 Gyrs (Bournaud & Combes 2002, Block et al 2002).

The influence of tidal interaction on the flatness of galaxy disks has been addressed by Reshetnikov et al. (1993) and Reshetnikov & Combes (1997) who found in interacting galaxies the ratio $h_z/h$ to be a factor of two higher, due to a disk thicker in absolute values, but also shorter in radial dimensions. Schwarzkopf & Dettmar (2000, 2001) confirmed this trend on a larger sample, but essentially because of a larger absolute thickness of interacting disks. They also noted that tidally interacting galaxies are more frequently warped and that the thickening was more noticeable in the outer parts.

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All these studies suggest that the relative flatness of a galaxy disk can vary by both effects during its evolution since the majority of galaxies have experienced interactions in the past. It appears that galaxies were dynamically “hotter” in the past (Abraham et al. 1999), which is seen in a smaller fraction of barred galaxies, and this could also have some consequences for the flattening. As the frequency of interactions/mergers strongly increases with redshift (e.g. Le Fèvre et al 2000), it is interesting to trace the evolution of the disk flatness with time by studying photometrically distant galaxies with high spatial resolution. This is now possible in the Hubble Deep Fields HDF-N and HDF-S and we present below the results of such a study. Section 2 describes our sample and Sect. 3 gives the derived general characteristics of disks in terms of radial and vertical profiles. The last section discusses the main results on the evolution of the flatness for stellar disks.

Throughout the paper, we adopt a flat cosmology with $\Omega_0 = 1$ and $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$. (For a model with $\Omega_m = 1/3$, $\Omega_\Lambda = 2/3$ and $\Omega_m + \Omega_\Lambda = 1$, linear sizes at $z = 0.7 - 1$ will be $\approx 25\%$ larger and absolute magnitudes are brighter by $\approx 0.5$.) If not differently specified, all magnitudes are expressed in the AB system (Oke 1974).

2. The sample

To study the structural parameters of distant edge-on spiral galaxies, we consider a sample of galaxies presented in Paper I (Reshetnikov et al. 2002). This sample includes 45 galaxies with a diameter larger than 1$'$ selected in the Hubble deep fields north and south. By applying a $V/V_{\text{max}}$ completeness test we derived that the sample is statistically complete – $V/V_{\text{max}}=0.50\pm0.04$ (Paper I). From the total sample of 45 distant galaxies we rejected objects with strong overlaps or interactions (e.g. n1031, SB-WF-2353-1914/2340-1898) and objects with unusual morphology (for instance, the "tailed" galaxies n16, n632, n1027). Also, we excluded the galaxies n749 and SB-WF-3329-3173 since they seem to be deviating from the edge-on orientation. The final sample consists of 34 edge-on galaxies suitable for further analysis.

The mean isophotal diameter of the galaxies within $\mu(I_{S14})=26.0$ is 2$''$3 (14 kpc at $z = 0.9$), the mean minor axis is $0''$.7. The latter is about 5 times larger than the FWHM of the PSF and therefore the galaxies are sufficiently resolved to study the vertical surface brightness distribution.

The general characteristics of the sample galaxies are presented in Table 1. Figure 8 (see Appendix) gives contour maps and radial major axis profiles of the objects.

3. General characteristics of edge-on disks

As judged from to the spectral energy distributions (Fernández-Soto et al. 1999), late-type galaxies (Sbc/Irr) predominate in our sample (Table 1). The mean redshift of the sample and the median of the distribution are the same, $z \approx 0.9$. The mean rest-frame absolute magnitude of the galaxies is relatively low $-M(B) \approx -18$th (Paper I). This is not unexpected since the Hubble deep fields have a very small coverage of only $\sim 2$ Mpc (comoving) at $z \sim 1$ each. At $z < 1$, the expected number of bright $L^*$ galaxies in the two fields is only $\sim 20-60$ (Ferguson et al. 2000). Therefore, assuming random orientation of galaxies, we can expect to find a few highly inclined bright spirals at $z < 1$.

Two of our edge-on galaxies are possible sources of sub-mm (SCUBA) radiation (Serjeant et al. 2002). These are the objects 304 and 774 in Table 1.

### Table 1. Edge-on galaxies in the HDF-N and HDF-S

| ID  | $I_{S14}$ | $b/a$ | $z$ | Type |
|-----|-----------|-------|-----|------|
| 15  | 24.70     | 0.28  | 0.45| Scd  |
| 55  | 23.32     | 0.20  | 0.18| Irr  |
| 120 | 24.27     | 0.24  | 1.60| Irr  |
| 153 | 24.94     | 0.35  | 0.73| Irr  |
| 273 | 22.47     | 0.45  | 0.905| Ell  |
| 304 | 25.60     | 0.35  | 1.78| Sbc  |
| 450 | 24.25     | 0.30  | 0.70| Sbc  |
| 476 | 22.77     | 0.39  | 0.421| Irr  |
| 506 | 24.47     | 0.25  | 0.751| Irr  |
| 534 | 23.82     | 0.25  | 0.95| Sbc  |
| 671 | 23.93     | 0.36  | 0.681| Irr  |
| 716 | 22.61     | 0.30  | 0.944| Sbc  |
| 727 | 23.93     | 0.33  | 0.904| Sbc  |
| 733 | 25.04     | 0.25  | 0.94| Irr  |
| 774 | 21.22     | 0.28  | 0.485| Sbc  |
| 805 | 25.33     | 0.30  | 0.66| Irr  |
| 817 | 22.19     | 0.40  | 0.485| Sbc  |
| 886 | 25.25     | 0.30  | 0.97| Irr  |
| 888 | 23.70     | 0.30  | 0.559| Irr  |
| 898 | 25.89     | 0.20  | 0.94| Irr  |
| 899 | 22.94     | 0.20  | 0.564| Sbc  |
| 938 | 23.03     | 0.27  | 0.557| Sbc  |
| 979 | 23.45     | 0.27  | 0.517| Sbc  |
| S0564 | 25.27    | 0.35  | 1.02| Irr  |
| S0566 | 25.67    | 0.30  | 2.16| Sbc  |
| S0578 | 23.16    | 0.22  | 0.98| Sbc  |
| S0747 | 25.31    | 0.24  | 0.98| Irr  |
| S1085 | 25.51    | 0.29  | 1.13| Irr  |
| S1404 | 23.09    | 0.20  | 0.50| Sbc  |
| S2661 | 24.10    | 0.30  | 0.92| Irr  |
| S2691 | 25.05    | 0.35  | 0.46| Irr  |
| S3053 | 26.07    | 0.30  | 1.05| Sbc  |
| S3458 | 23.40    | 0.40  | 0.69| Sbc  |
| S3685 | 26.63    | 0.35  | 2.23| SB1  |

The columns are: object identification (Paper I); $I_{S14}$ magnitude in the AB photometric system according to Fernández-Soto et al. (1999) (aperture magnitude corrected for the effect of neighboring objects and the flux lost in pixels outside the aperture); apparent axial ratio (from ellipse-fitting of the outer contours); redshift (spectroscopic redshifts are with three numbers, photometric redshifts obtained as the mean of redshifts presented by Fernández-Soto et al. 1999 and by Fontana et al. 2000); best-fit spectral type of galaxy (Fernández-Soto et al. 1999).
The mean observed axial ratio of the sample objects is 0.3 (Paper I, Table 1), a value exceeding the \( b/a \) ratio for local edge-on late-type galaxies. Possible reasons are a large intrinsic thickness of distant objects and also the influence of the point-spread function (PSF) on the apparent thickness of the small (\( \sim 2'' \)) HDF galaxies.

The largest uncertainty arises from the range of galaxy inclinations in our sample since we do not know the true intrinsic flatnesses of the sample objects and we also cannot use a standard approach, e.g., on the basis of the Hubble formula or from dust lane geometry. We selected our sample on the basis of careful visual inspection of the HDF frames. Secondary inclination indicators (asymmetry of the minor axis profiles, apparent shift of the nucleus location from the symmetry plane) are in the same range as in local samples of highly inclined spirals (e.g. de Grijs & van der Kruit 1996, Schwarzkopf & Dettmar 2000). Also, comparison with results of realistic numerical models for inclined spirals (e.g. Byun et al. 1994) indicates that most of our sample objects are within 5\(^\circ\) of edge-on orientation. Such a moderate deviation from edge-on orientation does not significantly change the slope of the vertical surface brightness distribution (van der Kruit & Searle 1981, Barteldrees & Dettmar 1994).

### 3.1. Optical colors

Figure 1 compares the observed colors of the 14 sample objects with spectroscopic redshifts and of nearly face-on spirals (randomly selected from the HDF-N) as a function of redshift. Both subsamples demonstrate the same trend of colors with \( z \) with a somewhat larger scatter for the edge-on disks. In the redshift interval \( z = 0.4 - 1 \) the edge-on disks show a hint of reddening in comparison with the face-on spirals: \( E_{u-b} = +0.25 \pm 0.18 \) and \( E_{b-v} = +0.23 \pm 0.13 \). These estimates correspond to inclination corrections for local spiral galaxies (e.g. de Vaucouleurs et al. 1991).

The predictions of apparent color evolution for two models of galaxy formation are shown in Fig. 1 according to Westera et al. (2002). The first model is characterized by a slowly growing dark halo with a continuous infall of gas and dark matter (accretion model, dashed line). The second one is a classical collapse model in a dark matter halo (collapse model, dotted line). Obviously, both models provide a satisfactory description of general observational trends of colors with redshift.

### 3.2. Major axis profiles

To investigate the global structure of the sample objects we adopted as a fitting model the simple standard double exponential disk: 

\[
L(r, z) = L_0 \exp(-\frac{r}{R_{\text{d}}}) - \frac{z}{h_z},
\]

where \( h \) and \( h_z \) are the scalelength and scaleheight, respectively. This simple model gives a reasonable description of the structure for most local disks, even in edge-on orientation (e.g. Misiriotis et al. 2000). Comparison with other vertical models for disks (for instance, with “\( \text{sech}^2(z/z_0) \)” or “\( \text{sech}(z/z_0^\prime) \)” distributions) can be done via \( h_z = z_0/\sqrt{2} \).

#### 3.2.1. Exponential profiles

Figure 8 presents the major axis photometric profiles of all galaxies in the \( I_{614} \) passband. The profiles show a large diversity in their shapes. Trying to fit the surface brightness distributions by an exponential law, we find that only 19 or 56\(\%\)\(\pm\)7\(\%\) of the disks are reasonably well described by exponential-like profiles.

The parameters for the exponential disks are presented in Table 2. Typical errors for the scalelength values are \( \approx 10\% \). The errors of scaleheights are presented in Table 2 and they are not just the formal errors of the least-squares fit, rather they were estimated by comparing re-
results from four independent fits (see Sect. 3.3). To check
the reliability of the exponential model to describe the
global photometric structure of the studied galaxies, we
have calculated total luminosities as \( L_{\text{exp}} = 2 \pi I_0 h^2 b/a \),
where \( I_0 = \text{dex}(-0.4m_0^\prime) \) (Table 2) and \( b/a \) is the appa-
rent axial ratio (Table 1). We found that the mean difference
of observed and model magnitudes is \( (I_{b14} - L_{\text{exp}}) = +0.31 \pm 0.12 \). This small difference can be attributed to the
infinite extrapolation for the model magnitudes.

The mean observational characteristics of the exponen-
tial disks are \( \mu_0(I_{b14}) = (22.99 \pm 1.16)^{\text{mag}}/\prime \), \( (h) = 0.43 \pm 0.11 \), \( (z) = 1.0 \pm 0.5 \). The observed values of \( \mu_0(I_{b14}) \)
(uncorrected for any inclination effects) have been con-
verted to a rest-frame \( B \) by applying the cosmological dim-
ming term and a \( k \)-correction color term for Irr galaxies
(Lilly et al. 1998, Lilly et al. 1995, Paper I). The resulting
mean characteristics of the central surface brightness and
of the scale length (after a small correction for the PSF –
see Sect.3.3) are \( \mu_0(B) = (20.7 \pm 1.3)^{\text{mag}}/\prime \), \( (h) = 2.3 \pm 0.7 \) kpc. For the disks with \( z \leq 1 \) (14 objects), the central
surface brightness is \( (21.05 \pm 1.0)^{\text{mag}}/\prime \). Therefore, dis-
tant edge-on disks demonstrate somewhat brighter central
surface brightnesses in comparison to nearby “Freeman”
disks for which \( \mu_0(B) = 21.5^{\text{mag}}/\prime \) in the AB sys-

tem (Freeman 1970). However, this difference is negli-
gible within the quoted scatter. Note also that our de-

erived value of the mean central surface brightness is in
good agreement with results for local samples of edge-
on galaxies: \( \mu_0(B) = (20.6 \pm 0.5)^{\text{mag}}/\prime \) (28 mem-
ers of interacting systems, Reshetnikov & Combes 1997),
\( (\mu_0(B)) = (20.7 \pm 1.1)^{\text{mag}}/\prime \) (45 normal galaxies, de Grijs
(1989).

As a first order approximation, let us assume that the
galaxy centers are optically thick, or close to that, and
so the central surface brightness does not depend strongly
on the inclination of the disk. This is a quite realistic ap-
proximation since at redshift \( z = 0.9 \) the filter \( I_{b14} \) corre-

sponds approximately to filter \( B \) (8140 \( \AA \):1+0.9)=4280 \( \AA \) and
the central regions of local spirals are not optically thin
in the blue spectral region. There are several lines of ar-
guments – discussed in the following – in favour of such a
conclusion.

The only direct and model-independent way to study
the transparency of the centers of galaxies is the chance
superposition of galaxies. In the innermost few hundred
parsec of NGC 3314A (a sub-L* Sc spiral galaxy) even
the most transparent regions between dust lanes show an
extinction of \( A_B \approx 7^m \) (Keel & White 2001). The disk of
NGC 3314A is inclined by \( i = 41^\circ \), so this galaxy is far
from edge-on orientation.

The same conclusion can be derived from statistical
arguments. For instance, Giovanelli et al. (1995), using a
sample of more than 1700 Sc galaxies with \( J \)-band pho-
metry, found only a weak dependence of the central sur-
face brightnesses of the disks on inclination: \( \mu_0(I) \) bright-
ens by about 0.3–0.4 mag between the face-on and edge-
on aspects. In the \( B \) band the dependence must be even
much weaker. Also, with regard to distant galaxies, Lilly et
al. (1998) did not find any significant correlation between
\( \mu_0(B) \) and disk axial ratio for spirals with \( z = 0.2 - 1 \).

Numerical simulations of the radiative transfer in
dusty disks show that even a small optical depth can sig-
nificantly reduce the expected projection effect on the cen-
tral surface brightness, especially in the \( B \) band (Byun et
al. 1994).

Therefore, assuming that the central surface bright-
ness shows only a weak (or negligible) dependence on the
galaxy inclination \( (\Delta \mu_0(B) = (0 - 0.5)^{\text{mag}}/\prime) \) we can
estimate an average “face-on” absolute magnitude of the
disks as \( M(B) = \langle \mu_0(B) \rangle - 5 \log(h) - 38.57 = -19.6^{\circ} \).
(\( h \) is the mean for the sample of 917 local late type
disks.)

The observational value for edge-on disks is \( -18^m \) fainter. The
value of dimming \( (1^{\circ}0\ldots1^{\circ}5) \) is consistent with uncertain-
ties with the value of global extinction for spirals (e.g.
de Vaucouleurs et al. 1991, Tully et al. 1998).

Table 2 presents also the \( h \) values in the \( V_{606} \) pass-
band. The mean scale length ratio is \( (h(I_{b14}))/h(V_{606})) = 1.04 \pm 0.13 \). (In the rest-frame this ratio resembles
\( h(B)/h(U) \).) Although we did not find any significant
color dependence as expected from the modeling by de
Jong (1996) for these filters, we cannot draw any firm
conclusions due to the small-number statistics.

In Fig. 2 we consider the \( \mu_0 - h \) plane. The dashed cross in
the figure shows the mean parameters of 917 Sb–Sd spi-
ral galaxies from Byun’s (1992) thesis (converted to \( B_{AB} \)).
The sample of Byun is angular diameter-limited (like ours)
and contains galaxies with angular diameter \( d \geq 1.7 \). The
mean absolute magnitude of this sample of local galaxies is \( M(I) = -20^m.4 \), or \( M(B) \approx -18^m.7 \) (adopting \( B - V = +1.7 \) for the disks of galaxies, de Jong 1996). Therefore, in edge-on orientation the local galaxies will show about the same luminosities as our sample galaxies. As one can see in the figure, the mean characteristics of nearby and distant sub-\( L^* \) disks are close (within quoted uncertainties) but \( z \sim 1 \) galaxies show a small displacement to higher values of surface brightness and smaller sizes. The shift to smaller values of \( h \) can be even larger due to the above mentioned effect of an apparent increase of scalelength in edge-on orientation.

At least three objects (n15, n805, and n898) can be considered as low surface brightness galaxies. These three galaxies show the rest-frame central surface brightness \( \mu_0(B) \approx 22 \) and are intrinsically faint \( (M(B) \geq -17^m) \). Adopting small or negligible internal absorption for such galaxies (e.g. Tully et al. 1998) and converting the observational values of \( \mu_0(B) \) to face-on orientation, we obtain \( \mu_0^f(B) \approx 23.5 - 24^m_{0.08}/\Omega'' \), typical for local low surface brightness galaxies (e.g. UGC 7321 – Matthews & Wood 2001).

3.2.2. Non-exponential profiles

15 of 34 disks (44\%±5\%) show “ridge”-like profiles, with several peaks and are often very asymmetric (see, e.g., n55, n506, or n886 in Fig. 8). Therefore, roughly half of \( z \sim 1 \) disks are not exponential. This fraction may be even larger due to preselection of the sample (Sect.2). This strongly contrasts with the local Universe and supports a late formation of disks (e.g. Mo et al. 1998, see van den Bergh et al. 2000 about galaxy morphology evolution with redshift).

**Fig. 3.** Reproduction of the HDF-N galaxy 866 (2."3×2."3) in the \( I_{814} \) passband.

To understand the origin of non-exponential brightness distributions, we have looked for possible counterparts for such objects among face-on galaxies in the deep fields. One of the best examples is the HDF-N galaxy 866 (2-153.0 in Williams et al. 1996) with the photometric redshift 0.96, \( I_{814} = 23.6 \) and the spectral type Scd (Fernández-Soto et al. 1999). The observed characteristics of this galaxy are close to those for the sample edge-on objects. This galaxy exhibits an irregular patchy disk with a low-contrast nucleus and embedded bright knots (five at least – see Fig. 3) and no spiral structure has developed yet. Figure 4 shows two photometric profiles of the galaxy which resemble the “ridge-like” profiles of many of our edge-on disks mentioned above. The galaxy is very similar in optical morphology to the peculiar galaxy H36490_1221 discussed in van den Bergh et al. (2001). Those authors suggest that such objects represent proto-late-type galaxies. Therefore, we can speculate that about half of our sample objects are young disks still in the formation process which will in their future evolution relax to an exponential distribution.

To check the frequency of non-exponential profiles among less distant spirals we selected in two deep fields all bright \( (I_{814} < 22^m) \) and non-edge-on galaxies of Sbc-Irr types. In total, we selected 17 galaxies with mean redshift \( \langle z \rangle \approx 0.5 \) and absolute magnitude \( \langle M(B) \rangle \approx -20^m \). Almost one third (5/17) of these objects show strongly non-exponential surface brightness distributions (work in preparation). The rest of the objects often demonstrate significant \( (\sim 0^m.5) \) deviations from the exponent. Therefore, non-exponential density distributions of galaxy disks are indeed very frequent at \( z > 0 \).

3.3. Minor axis profiles

The typical scalelengths of our objects are \( \approx 0'.3 - 0'.6 \) (Table 2). Therefore, \( h_z \) values are expected to be \( \leq 0'.1 \)

**Fig. 4.** Photometric profiles of HDF-N 866. Solid line – P.A.=125°, dashed line (the profile shifted down by 1") – P.A.=35°. The P.A. orientation is according to Fig. 3 (from top to the left).
Table 2. Characteristics of galaxies with exponential-like surface brightness distribution

| ID  | \( \mu_I^0 \) (" | \( h_I^0 \) (" | \( \mu_V^0 \) (" | \( h_V^0 \) (" | \( h_I^0/h_V^0 \) | \( h_I^0 \pm \sigma_{h_I} \) (kpc) | \( h_V^0 \) (kpc) |
|-----|----------------|---------|----------------|---------|----------------|----------------|---------|
| 15  | 22.94          | 0.38    | 23.27          | 0.46    | 0.18           | 1.70           | 0.31    |
| 153 | 23.17          | 0.29    | 23.69          | 0.28    | 0.21           | 1.50           | 0.32    |
| 304 | 23.70          | 0.32    | 24.05          | 0.35    | 0.27           | 1.73           | 0.25    |
| 450 | 23.34          | 0.48    | 23.79          | 0.48    | 0.33           | 2.62           | 0.34    |
| 476 | 21.67          | 0.40    | 22.22          | 0.46    | 0.27           | 1.74           | 0.09    |
| 671 | 22.54          | 0.39    | 23.25          | 0.37    | 0.23           | 2.03           | 0.16    |
| 727 | 22.28          | 0.36    | 22.50          | 0.34    | 0.18           | 1.98           | 0.06    |
| 733 | 23.76          | 0.59    | 24.54          | 0.70    | 0.17           | 3.32           | 0.11    |
| 744 | 20.42          | 0.31    | 21.57          | 0.32    | 0.19           | 1.40           | 0.03    |
| 805 | 23.56          | 0.34    | 24.03          | 0.35    | 0.14           | 1.74           | 0.07    |
| 817 | 20.46          | 0.47    | 21.55          | 0.47    | 0.4            | 2.26           | 0.21    |
| 888 | 22.52          | 0.37    | 23.03          | 0.35    | 0.24           | 1.80           | 0.07    |
| 898 | 24.24          | 0.47    |                |         | 0.17           | 2.66           | 0.09    |
| S0566| 24.00          | 0.42    | 24.43          | 0.52    | 0.18           | 2.24           | 0.10    |
| S0578| 22.36          | 0.64    |                |         | 0.25           | 3.70           | 0.20    |
| S0747| 23.85          | 0.57    | 24.36          | 0.56    | 0.13           | 3.16           | 0.08    |
| S1085| 23.69          | 0.51    | 24.00          | 0.41    | 0.10           | 2.76           | 0.03    |
| S3053| 24.38          | 0.57    | 25.4           | 0.7     | 0.11           | 3.12           | 0.14    |
| S3685| 24.00          | 0.28    | 23.86          | 0.24    | 0.22           | 1.43           | 0.11    |

The columns are: galaxy identification; central surface brightness and exponential scale length in arcsec (\( I_{814} \) passband); the same for the \( V_{606} \) filter; scale height to scale length ratio corrected for the PSF; final values of scale length and scale height in kpc.

It is important to note that \( h_z \) values are characteristics of the surface brightness gradients and are not linear sizes of any structures. To determine true scaleheights, we must study the effects of the HDF’s PSF on the surface brightness distribution of edge-on galaxies.

We have created a grid of double exponential disk models with various input parameters in the region \( h = 0.3’’ – 0.6’’ \) and \( h_z/h \) from 0.10 to 0.35. Then we prepared an empirical PSF by averaging several stellar images from the \( I_{814} \)-band frames of the HDF and convolved the model disks with the PSF. Comparing input and output disk characteristics, we determined empirical corrections for the observed parameters. The output parameters were estimated from the analysis of major and minor axes cuts of the model images. Finally, we found a strong influence of the PSF on the vertical scale of a disk. Also, we noted a small (within 10%) overestimation of the scale length values due to the PSF.

To present the influence of the PSF more clearly, we show in Fig. 5 the dependence of vertical stretching due to the PSF on the intrinsic flattening (\( h_z/h \)) of the model (the factor \( S \) gives the ratio of observed and true values of flattening). It is evident that intrinsically small (\( h \approx 0.3’’ \)) and thin (\( h_z/h \approx 0.1 \)) galaxies become more than 2 times thicker after application of the PSF. Therefore, we cannot expect to find very thin galaxies among small and distant HDF objects.

We excluded the central regions of the galaxies from the study in order to avoid any possible contribution of the bulges. (Although we do not see clear signs of bulges present in most galaxies – see Fig. 8). In order to explore the vertical structure of 19 exponential objects we examined one-dimensional cuts extracted at \( \pm (0.5–1)h \) in the galactocentric distance, avoiding the faintest outer regions of the objects. To minimize the influence of dust, we excluded the 1–2 central pixels of the cuts. In the analysis, we distinguished between the different sides of the galaxy.
planes and obtained the mean value of $h_z$ for each galaxy from 4 independent estimates. The mean observational (uncorrected for the PSF influence) scale height for the sample objects is $0.0713$, the mean dispersion $\sigma_{h_z} = 0.0703$. This value of dispersion (23%) characterizes the combined effect of the intrinsic variations of vertical scales and of errors in scale estimates.

The results of the analysis (corrected for the PSF) are given in the last three columns of Table 2. The mean relative thickness of the galaxies reduces after the PSF correction from $\langle h_z/h \rangle_{obs} = 0.31 \pm 0.09$ to $\langle h_z/h \rangle = 0.21 \pm 0.08$.

### 3.3.1. Vertical scale

The mean value of the vertical scale for distant edge-on galaxies is $\langle h_z \rangle = 0.46 \pm 0.21$ kpc or $\langle z_0 \rangle \approx 0.9$ kpc. The obtained mean value of the scale height is unbiased since we work with a statistically complete sample of galaxies with angular diameter larger than $1^\circ$ (Sect. 2, Paper I). These values are typical for normal spiral disks at $z \approx 0$ (e.g. Reshetnikov & Combes 1997, Schwarzkopf & Dettmar 2000). However, the direct comparison of local and distant samples must be carried out with caution due to very different selection effects. For instance, distant galaxies are sub-$L^*$ and relatively compact while the local samples usually contain brighter and larger disk galaxies (see Table 2 in Reshetnikov & Combes 1997).

In order to avoid (at least partially) the selection effects, we decided to compare our sample to local galaxies of the same size. Figure 2a shows the dependence of scaleheight on scalelength for 38 edge-on spiral galaxies ($T > 0$) in the $I$ passband according to de Grijs (1998). He adopted an exponential model for the observational brightness distribution and considered one-dimensional photometric profiles so we can directly compare our results. The mean regression for local galaxies predicts that a galaxy with $h = 2.3$ kpc must have $h_z = 0.3$ kpc. Therefore distant galaxies are thicker by a factor 1.5 on average. (In the non-zero $\Lambda$ cosmology with $\Omega_m = 1/3$, $\Omega_L = 2/3$ the relative thickening stays the same.)

Figure 2b presents the results of a two-dimensional decomposition in the $I$ passband of 34 edge-on spirals from the de Grijs (1998) sample recently presented by Kregel et al. (2002). The fitting model assumes a transparent disk (this is very questionable for our sample objects) and includes the line-of-sight integration of the luminosity distribution. Comparison of two methods shows that a one-dimensional method can overestimate the scalelengths by about 20% (Figure 3 in Kregel et al. 2002). However, a direct comparison of the $h$ values shows that $h$(de Grijs)/$h$(Kregel et al.) = 0.93 ± 0.20. A comparison of scaleheight values obtained by the two methods leads to a ratio $h_z$(de Grijs)/$h_z$(Kregel et al.) = 1.13 ± 0.22. Strictly speaking, both methods give a statistically insignificant difference. The Kregel et al. (2002) analysis results in a $\sim 1.4$ thickening of distant disks.

In Figure 3, we compare the Schwarzkopf & Dettmar (2000) sample of non-interacting galaxies (the data in the Johnson $R$ or Thuan & Gunn $r$ filters). The fitting model avoids the central dust lane in the line-of-sight integration. In this case the thickening factor increases to 1.9.

The last sample (Figure 4) includes 153 edge-on spiral galaxy with published photometric parameters in the $K_s$ band (Bizyaev & Mitronova 2002). The fitting model assumes a transparent disk. This work gives about a 1.8-fold thickening of distant galaxies.

#### 3.3.2. $h/h_z$ ratio

The relative thickness $h/h_z$ gives more reliable information about the possible thickening. For our sample we have $(h/h_z) = 5.4 \pm 2.0$, or, excluding one object for being most likely not exactly an edge-on galaxy (n817), 5.55 ± 1.9. Scaling this in terms of $z_0$, we obtain $(h/z_0) = 2.7 \pm 2.8$, with a median value of 2.63. The disks of local late-type galaxies are distributed over a wider range with the mean $h/z_0$ of $\approx 5$ (see Table 2 in Reshetnikov & Combes 1997). Moreover, disks of more late-type and gas-rich galaxies are thinner (Reshetnikov & Combes 1997, de Grijs 1998, Schwarzkopf & Dettmar 2000). For Sc/Sd disks Schwarzkopf & Dettmar (2000) give $h/z_0 \approx 8 \pm 2$. De Grijs (1998) and Kregel et al. (2002) give somewhat smaller ratios of 3.7 and 4.3, respectively. The Bizyaev & Mitronova (2002) sample shows the mean ratio of $h/z_0 = 4.8$.

In the previous discussion we neglected the possible dependence of scalelength on the passband. It is known (e.g. Table 7 in de Grijs 1998) that the $h$ value increases to the blue spectral region. The ratio of scalelength values in the $B$ and $K$ passbands can reach $\approx 1.5$, and $h(B)/h(I) \approx 1.3$ (de Grijs 1998). This effect, however, would only strengthen the thickness difference in the samples of local and $z \sim 1$ galaxies.

Therefore, $z \sim 1$ disks are definitely thicker in comparison to local galaxies. The factor for the vertical thickening ($\geq 1.5$) is approximately the same as for $z \approx 0$ interacting spiral galaxies ($\sim 1.5-2$: Reshetnikov & Combes 1997, Schwarzkopf & Dettmar 2000).

The dependence of $h/h_z$ values on the rest-frame central surface brightness is shown in Fig. 4. Thin disks are present for $\mu_{0}(B) \geq 21^{mag}/\arcsec^2$, while among bright objects ($\mu_{0}(B) < 21^{mag}/\arcsec^2$) we observe only thick disks with $h/z_0 \approx 2.2$. Taking into account that the $B$-band surface brightness characterizes the current rate of star formation, one can speculate that enhanced star-formation and enhanced thickness are the results of ongoing interactions and mergings. The observed morphology of galaxies does not contradict such a suggestion (Fig. 8).

### 3.4. Vertical scale for non-exponential galaxies

We examined the vertical profiles for 15 non-exponential galaxies at distances where the surface brightness drops by 1 mag relative to the nucleus. The mean observational value
Fig. 6. Scalelength – scaleheight relation for local spirals (open circles, (a) – de Grijs 1998, (b) – Kregel et al. 2002, (c) – Schwarzkopf & Dettmar 2000, (d) – Bizyaev & Mitronova 2002) and for our distant galaxies (solid rhombs). The solid lines show mean regressions for the nearby samples, the cross represents the average characteristics of distant galaxies and a $1\sigma$ scatter is indicated.

of the scaleheight obtained ($\langle h_z \rangle = 0.125$) is the same as for exponential objects (Sect.3.3) with a mean dispersion of individual values of 0.03. Taking into account that the redshift and angular diameter distributions in both subsamples are approximately the same, one can conclude that non-exponential disks must have approximately the same corrected value of the scale height ($h_z \sim 0.5$ kpc).

4. Discussion and conclusions

Our sample of distant edge-on galaxies shows an unexpectedly large fraction of non-exponential disks. From this observation we infer that many sub-$L^*$ disks are still in the formation process at $z \sim 1$, or the degree of tidal perturbations is higher, which has already been observed (e.g. Le Fèvre et al. 2000). One of the possible mechanisms to redistribute matter and to develop the exponential profile is the action of a long-lived bar (e.g. Combes et al. 1990). The young galaxies of our sample may not have experienced enough bar torques for this process, either because of very transient bars in very unstable disks, or because the effects of bars are continuously perturbed by mergers and accretion. The consequence is that a galaxy will spend much less of its time in a barred state. This conclusion is in good agreement with the general depletion of barred spirals at $z > 0.7$ (Abraham et al. 1999, van den Bergh 2002).

We have taken into account the effect of the PSF to demonstrate that the distant disks examined in this work are thicker by a factor of 1.5 if compared to local disk sample. This is a preliminary result based on a small number of sub-$L^*$ galaxies, and should be confirmed by larger statistical samples.

Since the star formation rate increases strongly with redshift (e.g. Madau et al. 1996, Flores et al. 1999, Thomson et al. 2001) one could speculate that the fraction of young stars formed in the midplane of disks is expected to be much higher in the past, making the young disk appear much thinner. Previous work on local samples has already concluded that tidal interactions between galaxies have a certain thickening effect of the order of 1.5-2 (Reshetnikov & Combes 1997, Schwarzkopf & Dettmar,
Fig. 7. Dependence of the relative thickness of distant galaxies on their rest-frame central surface brightness.

Since the frequency of galaxy interactions/mergers increases strongly with redshift, the present result fits into this picture. If one assumes that all disks have undergone such interactions, either additional cooling processes for the hot disks or a substantial growth predominantly in the radial direction are required to explain the flatness of disks in the local Universe.

However, other parameters should be taken into account in the interpretation. The distant disks are on absolute scales smaller than the local sample disks. It is likely that the morphological types are later (less evolved) and the masses of the galaxies smaller on average, as already found in previous studies based on HST surveys (e.g. Abraham et al. 1996, Glazebrook et al. 1998). The lower frequency of bars, associated with the smaller bulge-to-disk ratio, suggests that disks are more unstable and dynamically hotter (Abraham et al. 1999), which could explain the present thickening effect. In this context it is also noteworthy to mention that the galaxies of our sample contain only a few bulges as can be judged from the major axis profiles in Fig. 8. This can either be interpreted as evidence for a later formation of bulges by secular processes (e.g. Combes 2001) or that our sample is selected such that it represents early phases of late morphological types only.

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5. Appendix

We present here contour maps and major axis photometric profiles for the sample of edge-on galaxies in the HDF (Table 1).
Fig. 8. Contour plots (left) and major axis photometric profiles (right) of edge-on galaxies in the $I_{814}$ passband. The faintest contour corresponds to $2\sigma$ limiting isophote of the HDF or $\mu(I_{814}) \approx 25.75$, isophotes step is $0.''5$. Dashed lines – symmetric exponential model for the disk.
Fig. 8. (cont.)
Fig. 8. (cont.)