Allometry and Dissipation of Ecological Flow Networks

Jiang Zhang1, Lingfei Wu2*

1 School of Systems Science, Beijing Normal University, Beijing, China, 2 Department of Media and Communication, City University of Hong Kong, Hong Kong, China

Abstract

**Background:** An ecological flow network is a weighted directed graph in which the nodes are species, the edges are “who eats whom” relationships and the weights are rates of energy or nutrient transferred between species. Allometric scaling is a ubiquitous feature for flow systems such as river basins, vascular networks and food webs.

**Methodology:** The “ecological network analysis” can serve to reveal hidden allometries, the power law relationship between the throughflux and the indirect impact of node $i$, directly from the original flow networks without any need to cut edges in the network. The dissipation law, which is another significant scaling relationship between the energy dissipation (respiration) and the throughflow of any species, is also obtained from an analysis of the empirical flow networks. Interestingly, the exponents of the allometric law ($\eta$) and the dissipation law ($\gamma$) show a strong relationship for both empirical and simulated flow networks. The dissipation law exponent $\gamma$, rather than the topology of the network, is the most important factors that affect the allometric exponent $\eta$.

**Conclusions:** The exponent $\eta$ can be interpreted as the degree of centralization of the network, i.e., the concentration of impacts (direct and indirect influences on the entire network along all energy flow pathways) on hubs (the nodes with large throughflows). As a result, we find that as $\gamma$ increases, the relative energy loss of large nodes increases, $\eta$ decreases, i.e., the relative importance of large species decreases. Moreover, the entire flow network is more decentralized. Therefore, network flow structure (allometry) and thermodynamic constraints (dissipation) are linked.

Citation: Zhang J, Wu L (2013) Allometry and Dissipation of Ecological Flow Networks. PLoS ONE 8(9): e72525. doi:10.1371/journal.pone.0072525

Editor: Frederic Amblard, University Toulouse 1 Capitole, France

Received April 29, 2013; Accepted July 15, 2013; Published September 3, 2013

Copyright: © 2013 Zhang, Wu. This is an open-access article distributed under the terms of the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original author and source are credited.

Funding: National Natural Science Foundation of China under Grant No.61004107, www.nsfc.gov.cn. The funders had no role in study design, data collection and analysis, decision to publish, or preparation of the manuscript.

Competing Interests: The authors have declared that no competing interests exist.

* E-mail: zhangjiang@bnu.edu.cn

Introduction

An ecosystem is an open thermodynamic system driven by energy flows [1–3]. Each unit of an ecosystem is a living organism and is also an open system in which free energy is embodied by nutrient flows in the blood that are distributed from the heart to the cells through the vascular networks [4,5]. Moreover, all these units are connected by the so called food webs [6], on which energy flows are transported from environment to species and producers to consumers and are finally dissipated by respiration [2,3].

Allometric scaling is one of the most important and significant empirical laws for such open systems driven by flows [7–10]. Kleiber discovered a fundamental allometry between metabolism and body mass [7]. Subsequent studies further extended the allometric scaling laws to the community and ecosystem level and found that many characteristics of species, such as population density [11], abundance and trophic levels [12] in the food web, can also be predicted by body size.

The origins of the allometric scaling law can be explained by the fractal branching transportation networks occurring throughout living systems [4,5,13]. The most efficient transportation networks embedded in $d$-dimensional space, e.g., river basins in 2-d and vascular networks in 3-d, possess a ubiquitous power law, $C \propto A^\alpha$, where $A$ is the metabolism or input flows from the source to the network, $C$ is the total “mass” or the summation of all individual flow rates in the network, and $\eta = (d + 1)/d$ is the allometric exponent [5,14,15].

The food web, as the backbone of the ecosystem, can transport energy flow from the environment to each species [6,16–18]. Do food webs show the allometric scaling pattern? Garlaschelli gave a positive answer by reducing empirical food webs to spanning trees on which $A_i$ and $C_i$ defined as the number of nodes and the summation of the $A_i$ on the sub-tree rooted from $i$, respectively, have a significant power law relationship with ubiquitous exponents of approximately $\eta = 1.13$ for almost all the food webs collected for the purpose of the analysis [19–21]. Thus, spanning trees from food webs have structures similar to the transportation networks in living organisms. One of the highly salient features of Garlaschelli’s work is its focus on the network structure of food webs but not the individual or population of species as a whole as the metabolic theory usually does [9].

The energy flow information represented by the weights of the links in the network is, however, ignored by Garlaschelli’s work, although this information is already available in some empirical data sets [22]. Actually, the studies of ecological flow networks concerning both “who eats whom” binary relationships and the “in what rate” problem [23] have produced remarkable findings [2,24–26] and developed a sophisticated technique termed ecological network analysis [23,27–30].

Both allometric scaling laws in metabolic theory and ecological network analysis can contribute to the understanding of energy
flows. Nevertheless, they have hitherto remained separate. It appears that Garlaschelli’s approach should fill the gap between the two theories. However, this approach does not help due to its own limitations: many edges must be cut to obtain a spanning tree, and the flow information on the edges is never considered [19,31]. Therefore, it is crucial to extend Garlaschelli’s method to generalized flow structures.

In this paper, we extend Garlaschelli’s method by applying ecological network analysis to general flow networks through the calculation of the key variables $A_i$ and $C_i$ for each species in an ecological network. The allometric scaling relationships for 19 empirical ecological flow networks are presented in this study, and their allometric exponents are compared. In addition to the allometric scaling pattern, another interesting and significant empirical law called dissipation law is discussed. Dissipation, defined as the flux due to the respiration of the whole population, is found to be scaled with the total throughflow of each species on the ecological network. If we treat dissipation and total throughflow as the metabolic flux and body size of the whole population, the dissipation law can be viewed as an allometry at the population level.

In addition, a negative correlation between the allometric exponents ($\eta$) and dissipation law exponents ($\gamma$) is uncovered for the empirical networks collected. A method to perturb the flow distributions on original ecological flow networks is then developed to further investigate the scaling relationship linking $\eta$ and $\gamma$. The inverse relationship of $\eta$ and $\gamma$ is further confirmed through a large number of experiments on both empirical and artificial ecological networks. We conclude that $\gamma$ plays more important role than network topology in determining the allometric exponent.

This paper is organized as follows: Garlaschelli’s approach and its extensions to general flow networks are reviewed in Section Methods. The so called Flow Adjusting Algorithm, the perturbation method applied to ecological flow networks to study the relationship between $\gamma$ and $\eta$, is introduced in Subsection Flow Adjusting Algorithm. The allometric scaling rules and the dissipation laws for 19 empirical ecological flow networks are shown in Section Results. Also, the scaling relation between $\gamma$ and $\eta$ both on empirical and simulated flow networks is presented. Finally, we re-interpret the exponent in Section Discussion as the indicator of network centralization, i.e., the degree of the concentration of impacts on hubs in the network.

Methods

Review of Garlaschelli’s Method

To clarify the contribution of our method and its connection with current methods, we first review Garlaschelli’s method for a hypothetic food web. Figure 1(a), (c) shows how Garlaschelli’s approach can be applied to a hypothetical flow network (a). At first, a spanning tree (Figure 1 (a)) is constructed from the original network (Figure 1 (a)) by cutting edges. In this way, each sub-tree rooted from any $i$ can be viewed as a sub-system of the spanning tree. $A_i$ is the total number of nodes involved in this sub-tree and $C_i$ is the summation of the $A_j$ for each node in this sub-tree. Finally, the universal allometric scaling relationship of the $A_i$s and $C_i$s, with an exponent approximately 1.13, was found for all food webs considered [19].

Garlaschelli’s method was inspired by the Banavar’s model, which was developed to explain Kleiber’s law [5]. The spanning tree is simply Banavar’s optimal transportation network. Thus, energy flows into the entire system from the root along the links of the network to all the nodes. Suppose that each node consumes 1 unit of energy during each time step. A flux with 1 unit representing the energy consumption of each node should then be added to the original spanning tree (the dotted line in Figure 1[d]) [5]. As a result, the $A_i$ of each node is the throughput of this node. $C_i$ is the total throughflows of the sub-tree rooted from $i$. Essentially, the calculation of allometric scaling using Garlaschelli’s approach is based on this weighted flow network model [31].

Ecological Network Method

We will extend Garlaschelli’s method directly on the original weighted network without cutting edges (e.g. Figure 1[a]). However, the key question is how $A_i$ and $C_i$ are to be calculated for a general flow network.

According to the flow network interpretation of Garlaschelli’s method in Figure 1 (d), $A_i$ is the energy flux intake by node $i$, which is balanced with the total out flows through $i$. Thus, this concept can be extended to any flow network by defining $A_i$ as the throughput of node $i$ [27]. However, $C_i$ is not as straightforward to define as $A_i$ because the sub-system concept is not applicable if the considered network is not a tree.

To understand the meaning of $C_i$ in Garlaschelli’s method based on the flow network picture (Figure 1 (d)), we invoke the following hypothetical experiment. Assume that a large number of particles are flowing along the network (Figure 1 (d)), and that all particles passing any node, say $b$, will be attached a label, e.g., “b”. This trace marker will not be erased permanently unless the particle flows out of the network. We then find $C_b$ in Figure 1 (c) is simply the total number of particles labeled “b” that remain trapped in the entire network. This trace marker experiment can be also applied to other nodes independently and separately by attaching different labels so that all $C_b$s can be calculated according to the number of particles that ever passed node $i$.

This understanding can be extended to any flow network, whether or not it is a tree [31]. Although counting the number of labeled particles in the real network is difficult, we can perform this calculation directly with Markov chain techniques and the ecological network analysis developed by Patten et al [27,28]. As long as the flow network is in a steady state, so that all flows distributed on edges are stable, a fixed $C_b$ value can be calculated according to the flow structure.

Suppose that the flux matrix of the original network is $F$, in which each entry $f_{ij}$ stands for the flux from $i$ to $j$. Two special nodes, 0 and $S + 1$ ($S$ is the total number of species), representing the source and the sink, are contained in this matrix as the first(last) column(row). For most of the known ecological flow networks, the flows are balanced, i.e.,

$$\sum_{j=0}^{S} f_{ij} = \sum_{j=1}^{S+1} f_{ij}, \forall i \in [1,S]. \tag{1}$$

The webs not satisfying this condition will be balanced artificially by the method mentioned in Section A in File S1. Then, on a balanced flow network $F$, we define $A_i$ of $i$ as

$$A_i = \sum_{j=1}^{S+1} f_{ij}, \forall i \in [1,S] \tag{2}$$

which is the throughput of species $i$ in ecological network analysis [32].

To calculate $C_i$, we should convert the original flux matrix $F$ into a Markov chain $G$ in which each element is defined as...
$g_{ij} = f_{ij} / (\sum_{k=1}^{S} f_{ik})$ for all $1 \leq i, j \leq S$. Thus, $C_i$ can be calculated as,

$$C_i = \sum_{k=1}^{S} \sum_{j=1}^{S} f_{ij} \frac{n_{ji}n_{ik}}{n_{ij}},$$

where, $n_{ij}$ is an element of matrix $N$, termed the fundamental matrix [23,33] and defined as

$$N = (I - G)^{-1},$$

where, $I$ is the identity matrix. According to the ecological network analysis method, $C_i$ is simply the total number of particles that ever passed node $i$ (see Section B in the Supporting Information).

We then calculate $A_i$ and $C_i$ values for all nodes in the flow network to test the following allometric scaling law:

$$\text{Figure 1. Comparison of different methods for calculating the allometric scaling of a hypothetical ecological flow network.} (a) is a hypothetical network (the letter in each node is its index). The black node is the root, the numbers besides edges stand for flows, and the dashed lines represent dissipations; (b) shows the $A_i$ and $C_i$ values inside and beside node $i$, respectively, by our method; (c) is a spanning tree of (a), and the numbers inside and beside vertex $i$ are the $A_i$ and $C_i$ values calculated by Garlaschelli's method; (d) is the implicated flow network of (c), and the numbers beside the edges are the flux values.

doi:10.1371/journal.pone.0072525.g001

Allometry and Dissipation of Ecological Networks
where, $\eta$ is the allometric exponent that will be primarily discussed in the following sections.

Another significant scaling relationship found in this paper is the dissipation law,

$$D_i \propto A_i^\gamma,$$  \hspace{1cm} (6)

In this equation, $D_i$ is the flux from species $i$ to the sink in the original flux matrix before the imposition of the artificial balance. This value can be read directly from the original data as follows:

$$D_i = f_{iS+1}.$$  \hspace{1cm} (7)

In ecological flow network, this flux represents the respiration and output of species $i$. $\gamma$ is the dissipation exponent, which must be estimated from the data.

**Flow Adjusting Algorithm**

An important finding of this paper is the negative correlation between $\eta$ and $\gamma$. In addition to plotting the relationship between $\gamma$ and $\eta$ of the original flow networks directly, we also study the way in which the allometric scaling changes with the dissipation law by tuning the dissipation exponent without changing the topology on the networks. We devise a specific approach termed the Flow Adjusting Algorithm (FAA), to accomplish this goal.

Concretely, for a given flow network $F$, we keep the network topology and the ratio of influx ($f_{ij}/\sum f_{ij}$) unchanged, but adjust the flow distributions on edges to obtain a new flow network $F'$ such that:

1. The given dissipation law, i.e., $D_i \propto A_i^\gamma$, holds for every node $i$ in $F'$, where $\gamma$ is a given exponent which can be tuned. We will study how $\gamma$ impacts the allometric scaling law.
2. The flux balance condition, i.e., $\sum_j f'_{ij} = \sum_i f'_{ij}$, must be maintained for each node $i$.

The two requirements can be formalized by the following equations:

$$D'_i = c \left( \sum_{j=0}^{S} f'_0 \right) \gamma, \quad \forall i \in [1, S]$$  \hspace{1cm} (8)

where $c$ and $\gamma$ are given constants and the $f'_0$s and $D'_0$s are unknown variables. Therefore, we have a total of $2S$ equations but $E+S$ variables ($E$ is the total number of edges). For most flow networks, $E>S$, so infinite number of solutions to Equation 8 can be obtained. However, finding the solutions of these equations is still difficult due to their nonlinear nature.

The FAA is an approximate algorithm for solving these equations with the relative influx ratio of each edge unchanged. Suppose we have a node set $O$. Initially, we set time step $t=1$, $O^{(t)} = \{1,2,\ldots,S\}$ and the flux matrix as the original flow network $F^{(0)} = F$, where the superscript on $O$ and $F$ is the current time step. The algorithm will repeat the following steps:

1. For any node $i$ in $O^{(t)}$, the algorithm needs the current out flows from $i$, $\{f'_{ij} | f'_{ij} > 0\}$. Solve the following equation for $x_i$:

$$x_i = \sum_{j=1}^{S} f^{(t)}_{ij} + cx_i^\gamma,$$  \hspace{1cm} (9)

i.e., the new total influx to $i$. The solutions to these equations ensure that all nodes satisfy the flow balance condition and dissipation law.

2. Assign $x_i$ (the solution of Equation 9) to all incoming edges to $i$ proportionally, set

$$f_{ij}^{(t+1)} = x_i \frac{f_{0i}}{\sum f_{0i}^\gamma} \forall f_{ij} > 0$$  \hspace{1cm} (10)

and set,

$$D^{(t+1)}_i = c x_i^\gamma.$$  \hspace{1cm} (11)

Equation 10 guarantees that the ratio of every influx to the total influx remains unchanged.

3. Add all input nodes $j$s of $i$ into set $O^{(t+1)}$ and delete $i$ from $O^{(t+1)}$.

4. Set $t=t+1$ and repeat the previous steps until the stopping criteria, which include the requirements that the total running time is less than a given value and that the dissipation law exponent $\gamma$ and $R^2$ for the new flux matrix $F^{(t)}$ are close to the desired values, are satisfied.

This algorithm works only if the network is connected (i.e., there is at least one path from $0$ to every node $i$). For most networks, the algorithm can converge to a network satisfying the significant dissipation law (the R square of the power law regression between $D_i$ and $A_i$ is greater than a given value) with the given $\gamma$ exponent. However, it may oscillate on some topologies, especially for random networks (see Section D in the Supporting Information).

**Results**

**Description of Data Set**

Nineteen ecological flow networks for different habitats are considered in this study (Table 1). The networks are obtained from the online database (http://vlado.fmf.uni-lj.si/pub/networks/data/bio/foodweb/foodweb.htm). Most of the networks in this database are from the literature [25,34–39]. In Table 1, we list the name, the number of nodes $S$ and the number of edges $E$ of these networks. Note that the nodes are either living species or non-living compartments, and the weighted links are energy flows whose values vary substantially because the units and time scales of the measurements are very different. The source node $0$ and the sink node $S+1$ are the “input” node and the combination of “respiration” and “output” nodes in the original data, respectively. The dissipative flux of each node $i$ is simply the flow from $i$ to $S+1$ which can be read directly from the data. Most of the ecological flow networks considered are already balanced. The few unbalanced networks have been balanced with the approach cited in Section A of the Supporting Information.

**The Allometric Scaling Law**

We find that all the ecological flow networks possess a significant allometric scaling pattern. Their allometric exponents $\eta$ and the
values of $R_2^2$ are listed in Table 1. The values of $R_2^2$ are generally greater than 0.9 with the exception of the Rhode network, whose scale is relatively small ($S = 18$). All exponents $\eta$ fall within the interval $[0.83, 1.05]$. Most of these exponents are slightly greater than 1. An example of the allometric scaling fitting is shown in Figure 2.

To test whether the allometric scaling pattern is significant compared to random flow networks, we build a null model in which the numbers of nodes and edges are maintained, all links are re-connected randomly and all flows on edges are also randomly assigned on the interval $(0, f_{\text{max}})$ evenly, where $f_{\text{max}}$ is the maximum flux of the original network. From the inset of the upper plot in Figure 2, we see that the null model network does not show a significant allometric scaling law. We also attempt to compare the empirical flow networks against other null models. Some of these null models are obtained by maintaining the topology unchanged but randomly assigning the weights, whereas others are obtained by simply shuffling the weights among the edges. All the details of these null models are presented in Section C of File S1.

Among these null models, the models in which the weight information is retained generally have allometric scaling exponents similar to the original values. As a result, we know that it is the flow distribution, but not the topological structure of the network, that plays the more important role. However, it is impossible to study all possible flow distributions that can affect the final allometric law because hundreds of flows can be adjusted for an empirical flow network. Fortunately, we find that the dissipation pattern of flows provides a key clue for understanding the allometric scaling.

### The Dissipation Law

In ecology, the dissipative processes of a species have different forms, such as respiration, excretion, egestion, and natural and predatory mortality [3]. In our data, the dissipative flow is primarily respiration. It is plausible that this output flow would increase with the total throughflow of the focal species. However, it is not obvious for most collected ecological flow networks that the growth of dissipation for different species is slower than the

| Ecological Networks | Abbrev. | $S$ | $E$ | $\eta$ | $R_2^2$ | $\gamma$ | $R_c^2$ |
|---------------------|---------|-----|-----|-------|--------|--------|--------|
| Florida Bay, Dry Season | Baydry | 126 | 2102 | 1.010 | 0.995 | 0.915 | 0.949 |
| Florida Bay, Wet Season | Baywet | 126 | 2071 | 1.020 | 0.994 | 0.917 | 0.953 |
| Mangrove Estuary, Dry Season | Mangdry | 95 | 1462 | 1.010 | 0.997 | 0.978 | 0.983 |
| Everglades Graminoids, Dry Season | Gramdry | 67 | 863 | 1.030 | 0.999 | 0.973 | 0.997 |
| Everglades Graminoids, Wet Season | Gramwet | 67 | 863 | 1.020 | 0.999 | 0.977 | 0.998 |
| Cypress, Dry Season | CypDry | 69 | 639 | 0.998 | 0.996 | 0.957 | 0.949 |
| Cypress, Wet Season | CypWet | 69 | 630 | 0.997 | 0.997 | 0.965 | 0.988 |
| Mondego Estuary | - Zostrea site | Mondego | 44 | 401 | 1.010 | 0.999 | 0.979 | 0.997 |
| St. Marks River (Florida) | StMarks | 52 | 349 | 1.020 | 0.980 | 0.985 | 0.950 |
| Lake Michigan | Michigan | 37 | 210 | 1.010 | 0.999 | 0.995 | 0.999 |
| Narragansett Bay | Narragan | 33 | 194 | 1.010 | 0.991 | 0.813 | 0.942 |
| Upper Chesapeake | Bay in Summer | ChesUp | 35 | 203 | 1.050 | 0.997 | 0.952 | 0.991 |
| Middle Chesapeake | Bay in Summer | ChesMiddle | 35 | 195 | 1.040 | 0.996 | 0.851 | 0.761 |
| Chesapeake Bay | Mesohaline Net | Chesapeake | 37 | 160 | 0.994 | 0.997 | 0.985 | 0.985 |
| Lower Chesapeake | Bay in Summer | ChesLower | 35 | 163 | 1.050 | 0.997 | 0.926 | 0.971 |
| Crystal River Creek | (Control) | CrystalC | 22 | 107 | 1.040 | 0.997 | 0.959 | 0.995 |
| Crystal River Creek | (Delta Temp) | CrystalD | 22 | 83 | 1.040 | 0.998 | 0.963 | 0.996 |
| Charca de Maspalomas | Maspalomas | 22 | 82 | 0.956 | 0.966 | 1.150 | 0.737 |
| Rhode River Watershed | - Water Budget | Rhode | 18 | 54 | 0.828 | 0.866 | 1.200 | 0.963 |

doi:10.1371/journal.pone.0072525.t001

doi:10.1371/journal.pone.0072525.g002
growth of throughflow. A sub-linear relationship (Equation 6) between the dissipation and throughflow of each species holds per se. The values of the γs for different networks are also listed in Table 1. Note that almost all the γ values are less than 1 with the exception of the Maspalomas and Rhode networks. The lower plot in Figure 2 shows the dissipation scaling law for the Mondego flow network as an example.

The Relationship between γ and η
Because both γ and η are indicators for the entire flow network, to determine how these two numbers are correlated with each other we can simply plot different ηs against γs across all the collected empirical flow networks (see the blue dotted line in Figure 3). Although a general tendency for η to decrease with γ is observable, this trend is not significant for three major reasons. First, the number of sample points is too small to show a clear relationship. Second, most of the data points are concentrated in the circled area because all γs and ηs are similar in value for all networks. Finally, the original exponents are noisy, so that the η values fluctuate within the circled area.

To avoid the aforementioned problems, we investigate the relationship between γ and η by perturbing the original networks using the FAA cited in Subsection Flow Adjusting Algorithm. In this way, we can perturb the flow structure of the original network to obtain the expected dissipation law and to observe how exponent γ affects exponent η. Figure 4 shows the dependence of η on the perturbed networks according to the FAA based on the Mondego flow network, a randomized Mondego flow network (based on the Mondego’s topology but assigning flows randomly) and networks generated by the Niche model [40]. Note that the allometric scaling exponent decreases with the dissipation law exponent in a similar manner, whatever the original network structures. However, the shapes of the curves portraying the relationship between η and γ change with the network structures. For networks generated by the Niche model, η decreases less rapidly with γ if the connectances are higher (blue filled triangles versus purple hollow triangles and yellow filled diamonds vs. green hollow diamonds). As a result, the dissipation law exponent, but not the structure, is the major feature affecting the allometric exponent. However, we cannot conclude that the topological structure has no influence on the allometric exponents. Section D of File S1 will discuss the relationships among structure, the dissipation exponent and the allometric exponent in details for spanning trees.

In Figure 4, the red hollow circle in the middle of the black filled curve shows the combination of the dissipation law exponent γ of the empirical Mondego network F and the allometric exponent η of the adjusted flux matrix F’ resulting from the FAA. The red star corresponds to the original exponents of the Mondego network for both γ and η. We see that these two markers almost overlap. This result means that the original Mondego food web satisfies the following two requirements:

1) the dissipation law is significant for the original γ;
2) the balance equation, Equation 1, is obeyed exactly.

However, for other empirical networks, the perturbed result for the allometric exponent is not identical to the given original dissipation law exponent either because the flux balance requirement is violated or the dissipation scaling law is not significant (see Section E of File S1).

By adjusting the flows on all empirical flow networks with the fixed original dissipation law exponent γ, we can eliminate noise in the raw data. Because we require the networks to be flow balanced and satisfy the given dissipation law. The red solid curve in Figure 3 shows a clear relationship between γ and η for all empirical flow networks. Note that the decreasing trend of the

Figure 3. The relationship between γ and η in the original and adjusted ecological flow networks. The blue dashed line is the γ—η curve for the original flow networks, and the solid red line is the corresponding curve for the adjusted networks based on the same network structure and dissipation law according to the flow adjusting algorithm (see the main text).
doi:10.1371/journal.pone.0072525.g003
empirical exponents (blue dash line) is not as strong as that of the
exponents obtained from FAA (red line). Therefore, the two
requirements stated above are not satisfied perfectly by many
empirical ecological networks.

Discussion

Transportation Efficiency or Degree of Centralization?

Previous studies have explained the allometric scaling exponent
\( \eta \) as the transportation efficiency of the network, with values
ranging from 1 (a star network, the most efficient tree) to 2 (a
chain, the most inefficient tree). However, in our study, the
allometric scaling exponent is not restricted to the interval \([1,2]\).
Therefore, we should offer a new explanation for the exponent \( \eta \).
The key problem is to understand the indicator \( C_i \).

In Banavar’s model and Garlaschelli’s method, \( C_i \) is charac-
terized as a cost of energy transportation for the sub-tree rooted
from \( i \) [5]. From this perspective, the energy flows on the
redundant links (loops or cross-leveled links), except for the links in
the spanning tree, are wasted. Nevertheless, this interpretation
cannot be plausibly generalize to flow networks because (1) the
wasted energy in the weighted networks can be measured as the
dissipation of each node but not the weight of the edges and (2) all
energy links should be considered because they all contribute to
the overall flow distribution.

According to the particle tagging experiment described in
Section Methods, we can interpret \( C_i \) as the total impact of \( i \) to
the whole network along all flow pathways [31,41] because it is the
total number of particles passing \( i \) at least once. Thus, as
\( C_i \) increases, additional nodes will be affected by the particles
ever passed \( i \), and the direct and indirect influences of \( i \) will
increase. This interpretation of \( C_i \) can be extended to any flow
network.

As a species ascends up along the energy throughflow gradient
\( A_i \), its total impact \( C_i \) also increases with the relative speed \( \eta \)
according to the allometric scaling law \( C_i \propto A_i^\eta \). Therefore, the
important nodes (with larger \( A_i \)) may have much greater power
(total impact \( C_i \)) in the networks with larger exponents \( \eta \) than in
those networks with smaller exponents. For example, consider two
networks with four nodes. They have the same throughflow
distributions, e.g., \( A_1 = \{1,2,3,4\} \), but different exponents \( \eta_1 = 1 \)
and \( \eta_2 = 2 \). As a result, they have different \( C_i \) distributions,
\( C_i^{(1)} = \{1,2,3,4\} \) and \( C_i^{(2)} = \{1,4,9,16\} \). The most important node
\( (A_4 = 4) \) in the second network has a much greater impact
\( (C_i^{(2)} = 16) \) on the entire network than is the case in the first
network \( (C_i^{(1)} = 4) \). Hence, the concentration of the species impact
of the second network is larger than that of the first network. Thus,
the second network is more centralized than the first.

In short, the allometric exponent \( \eta \) measures the degree of
centralization of the network. The flow networks with exponent
\( \eta > 1 \) are centralized, whereas those with exponent \( \eta < 1 \) are
decentralized.

This new interpretation is compatible with the previous one. For
Garlaschelli’s spanning trees, the most centralized tree with a given
root is a chain but the most decentralized structure is a star (see the
discussion on spanning trees in Section D of the Supporting
Information).
Dissipation and Centralization

According to the new interpretation of the allometric exponent and the discussion of the relationship between dissipation and allometry, we can obtain an overall perspective: the networks may become more decentralized by dissipating more energy on larger nodes because the impacts of high flux nodes are weakened. In the networks with $\gamma > 1$, the dissipation flow per unit of throughflow increases with the throughflow, therefore the energy invested in the entire network decreases with the size of the node and the network is more decentralized. In contrast, if $\gamma$ is less than 1, the dissipating flux scales to the throughflow with a smaller relative speed. So that the large nodes can output more energy to the entire network to obtain a much more powerful impact on other nodes, the networks are more centralized.

However, an interesting and still unexplained finding is that the allometric exponents of empirical ecological networks are all close to 1. They are neither centralized nor decentralized. We conjecture that this finding can be explained by an optimization between flow structure stability and energy transport efficiency. This hypothesis will require additional study.

Concluding Remarks

In summary, this paper generalizes the universal allometric law to ecological flow networks and shows that the major factor influencing the allometric exponent is $\gamma$, the dissipation law exponent. By reinterpreting the allometric exponent as the degree of centralization, we establish a connection between network structure and thermodynamic constraints. This connection is of substantial importance and deserves additional attention.

References

1. Odum H (1988) Self-organization, transformity, and information. Science 242: 1132–1139.
2. Odum H (1983) System Ecology. John Wiley & Sons Inc.
3. Straskraba M, Jorgensen SE, Patten BC, Petten CC (1999) Ecosystems emerging: 2. dissipation. Ecol Model 117: 3–39.
4. West G, Brown JH, Enquist BJ (1997) A general model for the origin of allometric scaling laws in biology. Science 276: 122–126.
5. Banavar J, Maritan A, Rinaldo A (1999) Size and form in efficient transportation networks. Nature 399: 130–132.
6. Pimm S (2002) Food webs. University Of Chicago Press.
7. Kleiber M (1932) Body size and metabolism. Hilgardia 6: 315–353.
8. Brown J, West G (2000) Scaling in Biology. Oxford University Press.
9. Brown J (2004) Toward a metabolic theory of ecology. Ecology 85: 1771–1789.
10. West GB, Brown JH (2005) Review the origin of allometric scaling laws in biology. Ecol Model 117: 3–39.
11. Enquist BJ, Brown JH, West GB (1998) Allometric scaling of plant energetics and population density. Nature 395: 163–165.
12. Hechinger RF, Lafferty KD, Dobson AP, Brown JH, Kuris AM (2011) A common scaling rule for abundance, energetics, and production of parasitic and free-living species. Science 333: 445–448.
13. West G, Brown J, Enquist B (1999) The fourth dimension of life: Fractal geometry and allometric scaling of organisms. Science 284: 1677–1679.
14. Rodriguez-Inurribi I, Rinaldo A (1997) Fractal River Basins. Chance and Self-Organization. Cambridge: Cambridge University Press.
15. Dreyer O (2001) Allometric scaling and central source systems. Phys Rev Lett 87: 038101.
16. Cohen J, Briand F, Newman C (1990) Community Food Webs: Data and Theory. Biomathematics Vol. 20. Berlin: Springer.
17. Pascual M, Dunne JA (2003) Ecological Networks: Linking Structure to Dynamics in Food Webs. Oxford University Press.
18. Allestria S, Bodini A (2004) Who dominates whom in the ecosystem? energy flow bottlenecks and cascading extinctions. J Theor Bio 230: 351–358.
19. Garlaschelli D, Caldarelli G, Pietronero L (2003) Universal scaling relations in food webs. Nature 423: 165–168.
20. Camacho J, Arenas A (2003) Food-web topology: Universal scaling in food-web structure? Nature 437: E3-E4.
21. Gaba F, Marrell D (2003) A simple explanation for universal scaling relations in food webs. Ecology 86: 3258–3263.

Supporting Information

File S1 Contains. Figure S1. The relationships between $\gamma$ and $\eta$ for different balancing methods. Naive, input, output, and average stand for different balancing methods. Figure S2. An example network showing the calculations of $A_i$ and $C_i$. Figure S3. Allometric Scaling Patterns forNull Models based on the Mondego Network. Figure S4. Comparison of Allometric Scaling Exponents of Original Networks with the Null Models. Figure S5. Dissipation Law for Null Models of Mondego Network. Figure S6. $\eta$ and $R^2$ Change with $\theta$, $\beta$, and $\gamma$. Figure S7. $\eta$ and $R^2$ Change with $\gamma$ for different random networks. Figure S8. $\eta$ and $R^2$ Change with $\gamma$ and $\beta$ on the Random Networks based on Spanning Trees. Figure S9. $\gamma$ and $\eta$ Relation Adjusted by the Flow Adjusting Algorithm on the Collected Ecological Networks. Figure S10. Dissipation Scaling Law for the Original and Balanced Flow Networks of Baydtry and Rhode. Table S1. Unbalanced Flux for Each Ecological Network. Table S2. Comparison among Different Balancing Methods for Allometric Scaling Exponent. (PDF)

Acknowledgments

We acknowledge the discussion with Prof. Dunne in Santa Fe Institute.

Author Contributions

Conceived and designed the experiments: JZ. Performed the experiments: JZ. Analyzed the data: JZ LW. Contributed reagents/materials/analysis tools: JZ LW. Wrote the paper: JZ LW.

Allometry and Dissipation of Ecological Networks