X-ray extinction from interstellar dust
Prospects of observing carbon, sulfur and other trace elements

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ABSTRACT

Aims. We present a study on the prospects of observing carbon, sulfur, and other lower abundance elements (namely Al, Ca, Ti and Ni) present in the interstellar medium using future X-ray instruments. We focus in particular on the detection and characterization of interstellar dust along the lines of sight.

Methods. We compare the simulated data with different sets of dust aggregates, either obtained from past literature or measured by us using the SOLEIL-LUCIA synchrotron beamline. Extinction by interstellar grains induces modulations of a given photoelectric edge, which can be in principle traced back to the chemistry of the absorbing grains. We simulated data of instruments with characteristics of resolution and sensitivity of the current Athena, XRISM and Arcus concepts.

Results. In the relatively near future, the depletion and abundances of the elements under study will be determined with confidence. In the case of carbon and sulfur, the characterization of the chemistry of the absorbing dust will be also determined, depending on the dominant compound. For aluminum and calcium, despite the large depletion in the interstellar medium and the prominent dust absorption, in many cases the edge feature may not be changing significantly with the change of chemistry in the Al- or Ca-bearing compounds. The extinction signature of large grains may be detected and modeled, allowing a test on different grain size distributions for these elements. The low cosmic abundance of Ti and Ni will not allow us a detailed study of the edge features.

Key words. ISM: dust, extinction – X-rays: ISM – X-rays: individuals: GX 5-1, GX 340+00, GX 3+1 – techniques: spectroscopic

1. Introduction

Absorption and scattering in the X-ray band has proved a useful diagnostic of the interstellar dust (ID) properties. By virtue of the broad band coverage, the X-ray band displays many photoelectric absorption edges, caused by the mixture of gas and dust intervening along the line of sight towards bright background sources (Draine 2003; Hoffman & Draine 2016). Absorption by interstellar grains is detected as a result of the interaction between the incoming X-ray photon and the electrons inside the grain’s atoms. The multiple-generated photoelectron-waves interfere with each other both constructively and destructively. This interference pattern depends on the complexity of the chemical compound and the distance of the electrons from the nucleus. Each pattern is a fingerprint of a given material (Rehr & Albers 2000). The extinction cross section, the sum of the absorption and scattering cross section (e.g. Draine 2003; Corrales et al. 2016) in the X-ray band, provides, in principle, not only direct estimate on the chemistry of the interstellar medium (ISM), but also information on the size distribution, crystallinity and porosity of the intervening grains (Hoffman & Draine 2016; Zeegers et al. 2017; Rogantini et al. 2018).

Early studies already pointed out that absorption by the ISM contributed to the shape of X-ray spectra (Schattenburg, & Canizares 1986; Paerels et al. 2001; Juett et al. 2004). However, in recent years, the deep features of the Fe L and O K and Si K edges have been recognized to be largely caused by dust absorption (e.g. Lee et al. 2002; Ueda et al. 2005; Schulz et al. 2016) and have been studied using the grating spectrometers on board the X-ray Observatories Chandra and XMM-Newton. These studies made use of absorption profiles either taken from the literature (Costantini et al. 2012; Pinto et al. 2016, 2017; Valencic & Smith 2013) or obtained with dedicated synchrotron measurements (Lee et al. 2009; Zeegers et al. 2017).

Outside the energy band where the sensitivity and resolution of the current instruments is maximized, it is at this moment challenging to study interstellar dust. An example is given by the tentative study of the C K edge (Schneider & Schmidt 2010), which was severely hampered by various instrumental effects, although the carbon edge would formally be included in the energy range of Chandra-LETGS.

In this paper we investigate the prospective of observing and modelling the elements of the ID which have not been yet studied, which happen to fall in the 1.5–8.3 keV band (Al, S, Ca, Ti, Ni) and at E<0.5 keV (carbon). For discussion on how the features that can be currently studied (O K, Fe L Mg K, and Si K...
Fig. 1. Upper panel: abundance pattern as a function of energy for the absorbing elements in the X-ray band. The K-edge energy is indicated, except from Fe, for which both the K- and the L-edges, at 7.1 and \( \sim 0.7 \) keV, respectively, can be studied. Abundances follow Lodders (2010) and they are expressed in terms of \( \log (X/H) + 12 \). In this frame, the abundance of hydrogen is 12. The open diamond mark the elements that are accessible by current instruments. The triangles are the relevant elements that will be accessible by future instruments to study dust. The black triangles are the subject of this work. Middle panel: range of depletions as reported by Jenkins (2009) for all elements, except: C (Jenkins 2009; Whittet 2003), F (Snow et al. 2007), Na (Turner 1991), S (Gry & Jenkins 2017), K (Snow 1975), Ca (Crinklaw et al. 1994), Co (Federman et al. 1993), Al (Jenkins & Wallerstein 1999), Ar (Soña & Jenkins 1998). Lower panel: energy range covered by present and future (red) mission. The solid line highlights the energy range where the instrument capabilities are optimal for observing absorption by dust.

In Fig. [1] we show the abundance pattern of the photoelectric absorption edges of the elements (with atomic number \( A = 6-30 \)) as a function of the X-ray energy. The empty diamonds mark the edges that have been already studied by current instruments unveiling the dust features: the O K and Fe L edges at 0.534 and 0.7 keV respectively (Lee et al. 2009; Costantini et al. 2012; Pinto et al. 2010, 2013; Valencia & Smith 2013); the Mg and Si-K edges at 1.3 and 1.84 keV respectively (Zeegers et al. 2017, 2019, Rogantini et al. in prep.). The black triangles mark the edges presented in this work. We present the Fe K edge, marked with a light gray triangle in the figure, in a separate paper (Rogantini et al. 2018), but see also Lee & Ravel (2005). The middle panel shows the range of depletion, defined as the amount of dust over the total amount of matter in the ISM (dust and gas) that is expected for a given element. The wide range of depletions for some elements is due to the different density environments where those are observed (e.g. Jenkins 2009).

In the lower panel of Fig. [1] we show the energy range of present and future missions. The solid line highlights the region where the instrument capabilities (resolution and effective area) are optimal to observe the dust absorption features. The Chandra and XMM-Newton observatory (both launched in 1999, Weisskopf 1999, Jansen et al. 2001) are still in operation. The grating spectrometer Arcus will cover the soft X-ray range (Table 1). It is a NASA mission currently in the study phase (Smith et al. 2016). The calorimeter on board the X-ray Imaging Spectroscopy Mission (XRISM, to be launched around 2021) is planned to have the same characteristics of the one on board of the lost Hitomi satellite (Mitsuda et al. 2014). Finally, we display the energy coverage of the Athena calorimeter XIFU (Barret et al. 2016), to be launched in 2030. Both XRISM and Athena’s calorimeter will be optimal to observe the higher energy dust features (Table 1).

1.1. The elements in this study

One of the major players in the ID, carbon, constitutes around 20% of the total depleted mass in the Galaxy (Whittet 2003). Its depletion covers a relatively narrow range of values, (Figure [1] showing that it is not a strong function of environmental density. It has been hypothesized that the majority of carbon should be locked in graphite grains, as a likely explanation for the 2175 Å emission feature (Draine 1989, 2003, and references therein). Graphite has been commonly adopted in ID mod-

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The energy resolution is 2.5 and 5 eV for XIFU and Resolve, respectively. Arcus resolution is defined by its resolving power R ~ 3000 over the 0.5–1 keV band.

Sulfur in dust has been also detected near C-rich AGB stars, planetary nebulae (Hony et al. 2002) and protoplanetary disks (Keller et al. 2002), predominantly in form of troilite (FeS). Finally, sulfur is abundant in solid form in planetary systems bodies, such as interplanetary dust particles, meteorites and comets (e.g. Wooden 2008 and references therein).

The presence of sulfur in dust form in the ISM has been suggested in association with GEMS (Glasses with Embedded Metal and Sulfides, Bradley 1994), where the FeS particles would be more concentrated on the surface of the glassy silicate. However, the majority of GEMS may well be of nebular origin, rather than the ISM (Keller & Messenger 2008). Sulfur in FeS, consistent to be of ISM origin, has been recorded in the data from the Stardust mission (Westphal et al. 2014). This evidence revitalizes the idea of the presence of sulfur in dust-form also in less dense environments of the ISM. The presence of strong UV radiation and cosmic rays has been thought to be the cause for the extreme sputtering of the highly volatile S, for example from GEMS surfaces. Recent experiments however put to the test this hypothesis (Keller et al. 2010), showing that UV bombardment has in fact little influence on sulfur stuck on a grain surface.

Aluminum, calcium and titanium are extremely depleted in the ISM (Fig. 11). Ca and Ti show also a very similar depletion pattern as a function of gas density (Crinklaw et al. 1994). These two elements are found in gas only in tenuous environments, associated with warm inter cloud media in both the halo (Edgar & Savage 1989) and the disk of the Galaxy (Crinklaw et al. 1994). The depletion of Ti is severe, regardless of the environment, ranging between -1 dex and -3.1 dex (Wetlty & Crowther 2010). The ratio of column density between Ti and Ca, representative of the element neutral gas phase, is in general constant in the Galaxy (~0.4, Hunter et al. 2006). Al, Ca and Ti have a similar condensation temperature (1400–1600 K, Field 1974) and it has been hypothesized that, being the first to form in e.g. a stellar envelope or a supernova environment, they would form the core of complex dust grains with silicate and possibly ice mantles (e.g. Clayton 1978). This would provide a natural shield for these Al, Ca, Ti-bearing compounds, preventing their destruction and ensuring a high depletion in the vast majority of the environments. Under the condition of thermodynamic equilibrium, aluminum first condenses in Al2O3. From there it may evolve into spinel (MgAl2O4) and eventually into a Ca and Al-bearing silicate. The latter are stable compounds, thanks to very high binding energies (Trivedi & Larimer 1981). Calcium is mostly locked in dust silicates (e.g. CaMgSi2O6, Field 1974; Trivedi & Larimer 1981). Calcium carbones, possibly formed in AGB stars envelopes (e.g. Kemper et al. 2002), are believed to be unstable and therefore Ca inclusion in silicates, which form already at high temperatures, are favored (Ferrarotti & Gail 2005).

Titanium is produced by AGB stars mostly in the form of TiO2, which constitute a seed nucleus later included in the larger/coated grains (e.g. Ferrarotti & Gail 2006).

Nickel depletion is found to correlate with the one of Fe for a variety of environments, from planetary nebulae (Delgado-Inglada et al. 2016) to diffuse interstellar clouds (Jenkins 2009). These two elements display a similar condensation temperature (1336 and 1354 K for Fe and Ni, respectively, Wesson 1989), which already point to a simultaneous inclusion

Table 1. Parameters of the instruments used in the simulations at the energy of the elements studied here.

| Element | R/Å | $\Delta E_{\gamma}$/cm$^{-1}$ | Inst.
|--------|-----|-----------------|------
| C      | 2540 | 369             | Arcus
| Al     | 621  | 11022           | XIFU
| S      | 986(552) | 5949(209) | XIFU(Resolve)
| Ca     | 1611(575) | 3734(271) | XIFU(Resolve)
| Ti     | 1984  | 2742            | XIFU
| Ni     | 3331  | 704             | XIFU

Notes: the energy resolution is 2.5 and 5 eV for XIFU and Resolve, respectively. Arcus resolution is defined by its resolving power R ~ 3000 over the 0.5–1 keV band.
Rogantini et al. (2018) to obtain the extinction profiles to be confronted to the astronomical data. The laboratory data are trans-formed to transmission spectra and matched (via χ²-fitting) to tabulated transmission data of the same compound, where we assume an optically thin sample, as to mimic the conditions in the ISM. In doing this we only fit the pre- and post-edge of the tabulated data, leaving the edge energy as measured in the laboratory. The transmission tables are provided by the Center for X-ray Optics at Lawrence Berkeley National laboratory based on tabulated data by Henke et al. (1993). From the transmission spectra, the attenuation coefficient can be obtained and consequently we can obtain the imaginary part of the refractive index from this coefficient. The real part of the refractive index m, is calculated via the Kramers-Kroning relations (Bohren 2010). The knowledge of m is needed in order to calculate the extinction cross section to involve both the effect of absorption and scattering. The extinction cross sections are calculated using Mie theory (Mie 1908, Wiscombe 1980) for C, Al and S. We used instead the anomalous diffraction theory (ADT, Van der Hulst 1957) for Ca, Ti and Ni. The ADT theory can be used when the ratio x = 2πa/λ ≫ 1 where a is the grain size and λ is the wavelength of the incident radiation. In order to obtain the extinction cross section for a range of dust grain radii, we assume the MRN size distribution (Mathis et al. 1977 and Sect. 3). Once the absolute cross section as a function of energy has been obtained, we implemented the extinction profiles in the already existing AMOL model in the fitting code SPEX (ver. 3.03, Kastra et al 1996). The AMOL model is an absorption model to be applied to the emitting continuum model of a source. It allows to fit the X-ray edges for the column densities of a set of four dust compounds at a time (see also Zeegers et al. 2017). In real astronomical observations, the dust extinction feature will always coexist with the gas feature of the corresponding element. The absolute energy of the edges of the gas phase may be reported in the literature at different energies, with discrepancies sometimes of few eV. In SPEX, the gas edge energies, for the elements in this work, are implemented following Verner et al. (1996). In this paper, we apply a shift to the laboratory data in order to consistently compare them with the gas edge features as seen by SPEX. Discrepancies among different measurements and theoretical calculations may be found in the literature. High resolution X-ray spectroscopy will help determining the absolute energy scale of the edge (e.g. Gorczyca et al 2013).

2. Extinction profiles

In this paper, we make use of literature values to infer the absorption absolute cross sections for all elements, except for Al (Sect. 2.1). Measurements of X-ray edges profiles are mostly carried out for industry and are rarely of interest for astronomical applications. For this reason, the sample selection is bound to be incomplete. The compounds used are listed in Table 2. We follow closely the method presented in Zeegers et al. (2017) and Rogantini et al. (2018) to obtain the extinction profiles to be confronted to the astronomical data. The laboratory data are transformed to transmission spectra and matched (via χ²-fitting) to tabulated transmission data of the same compound, where we assume an optically thin sample, as to mimic the conditions in the ISM. In doing this we only fit the pre- and post-edge of the tabulated data, leaving the edge energy as measured in the laboratory. The transmission tables are provided by the Center for X-ray Optics at Lawrence Berkeley National laboratory based on tabulated data by Henke et al. (1993). From the transmission spectra, the attenuation coefficient can be obtained and consequently we can obtain the imaginary part of the refractive index from this coefficient. The real part of the refractive index m, is calculated via the Kramers-Kroning relations (Bohren 2010). The knowledge of m is needed in order to calculate the extinction cross section to involve both the effect of absorption and scattering. The extinction cross sections are calculated using Mie theory (Mie 1908, Wiscombe 1980) for C, Al and S. We used instead the anomalous diffraction theory (ADT, Van der Hulst 1957) for Ca, Ti and Ni. The ADT theory can be used when the ratio x = 2πa/λ ≫ 1 where a is the grain size and λ is the wavelength of the incident radiation. In order to obtain the extinction cross section for a range of dust grain radii, we assume the MRN size distribution (Mathis et al. 1977 and Sect. 3). Once the absolute cross section as a function of energy has been obtained, we implemented the extinction profiles in the already existing AMOL model in the fitting code SPEX (ver. 3.03, Kastra et al 1996). The AMOL model is an absorption model to be applied to the emitting continuum model of a source. It allows to fit the X-ray edges for the column densities of a set of four dust compounds at a time (see also Zeegers et al. 2017). In real astronomical observations, the dust extinction feature will always coexist with the gas feature of the corresponding element. The absolute energy of the edges of the gas phase may be reported in the literature at different energies, with discrepancies sometimes of few eV. In SPEX, the gas edge energies, for the elements in this work, are implemented following Verner et al. (1996). In this paper, we apply a shift to the laboratory data in order to consistently compare them with the gas edge features as seen by SPEX. Discrepancies among different measurements and theoretical calculations may be found in the literature. High resolution X-ray spectroscopy will help determining the absolute energy scale of the edge (e.g. Gorczyca et al 2013).

2.1. Laboratory data for Al

For Al, we made use of the laboratory data that we collected at the LUCIA beamline at the Soleil synchrotron facility which offers an energy resolution of ~0.25 eV. Both the samples, Al₂O₃ and MgAl₂O₄, were commercially available from the Alfa-Aesar and Aldrich company, respectively. The samples, in powder form, were pressed on thin indium foil, placed on a copper support which was placed in a vacuum environment. The sample was then irradiated by an X-ray beam of which the energy is tuneable. The X-ray fine structures were measured through fluorescence. At these soft X-ray energies, this method is more practical than the more intuitive method of measuring the transmission through the sample, because for transmission measurements the samples have to be too thin to be easily handled. The fluorescent method to obtain the XAFS does require a correction for possible saturation. This correction was performed with the program FLUO, which is part of the UWXAFS software (Stern et al 1995). A full description of the procedure for the analysis of the data can be found in Zeegers et al (2017).

3. Simulations

We present the prospects of detecting absorption K edges relevant for dust studies using future missions (Fig. 1). The only instrument proposed for studying the soft X-ray energy band is,
at this moment, the Arcus grating spectrometer (Smith et al. 2016). For the energy above \( \sim 2\) keV, two microcalorimeters will provide an unprecedented resolution: Resolve and XIFU, on board of XRISM (Smith et al. 2016) and Athena (Nandra et al. 2013), respectively. The effective area and the resolving power of these three instruments at the energy of the features studied here are reported in Table 1. With the chosen exposure time we would obtain an associated error on the dust or gas column density of around 1% in Table 1. With the chosen exposure time we would obtain an associated error on the dust or gas column density of around 1% for C (Arcus), Al, S and Ca (XIFU).

The simulations were carried out having in mind realistic scenarios in our Galaxy, in order to prove the effective prospect of future instruments to measure physical parameters. We first simulated the data, considering the different instruments responses and including noise, assuming first that the photoelectric edge is only due to gas absorption. These simple simulations are associated error on the dust or gas column density of around 1% for the XRD simulation, folded with the appropriate response, and the resolving power of these three instruments at the energy of the features studied here are reported in Table 1. With the chosen exposure time we would obtain an associated error on the dust or gas column density of around 1% for C (Arcus), Al, S and Ca (XIFU). The simulations were carried out having in mind realistic scenarios in our Galaxy, in order to prove the effective prospect of future instruments to measure physical parameters. We first simulated the data, considering the different instruments responses and including noise, assuming first that the photoelectric edge is only due to gas absorption. These simple simulations are confronted with models, folded with the appropriate response, which include an amount of gas set by the typical depletion found in the literature for a given element plus the contribution of a dust compound. We adopted the MRN dust size distribution in all cases (Mathis et al. 1977). The MRN model offers a simple parameterization of the dust size distribution: \( n(a) \propto a^{-3.5} \). It has been found to approximate at first order the conditions in our Galaxy for grain size between 0.025 and 0.25 \( \mu \)m. However, the exceptional depletion of Al and Ca, joint to favorable observing conditions allowed us to test also the detectability of a distribution where the mass distribution is skewed towards larger grains (Draine \\& Frasch 2009). This distribution has a size range of \( a = 0.02 \sim 1 \mu \)m, with an average grain size of 0.6 \( \mu \)m. The observing conditions are favorable for these two elements first because the brightness of the sources with a favorable \( N_H \) is often high (e.g. low mass X-ray binaries, LMXB, near the galactic center) and, second, the edges fall in a large effective-area region of the instrument. Unless otherwise stated, we chased the maximum depletion allowed by previous studies (Tab 1). This is a reasonable assumption in most cases. Indeed, for the photoelectric edge to be detected, a substantial column density is required. This also indicates a dense environment, where the depletion is large. Here below we describe the conditions under which the simulations were performed for each element.

Carbon: the depletion of carbon has been optimistically assumed to be 0.6, but still implies a substantial role in gas absorption in the carbon edge. For this reason, although the edge feature itself can also be detected at relatively small column density (e.g. \( N_H > 10^{20} \) cm\(^{-2}\)), a much larger amount of matter is necessary to make the dust features evident. Here we simulated a column density of \( N_H = 1.6 \times 10^{21} \) cm\(^{-2}\) for a flux in the soft energy band of \( 3 \times 10^{-9} \) erg cm\(^{-2}\) s\(^{-1}\) (Fig. 2). Note that for this column density, the C\(_\alpha\) absorption lines from gas are already saturated (Fig. 2), therefore they cannot be used straightforwardly to measure depletion. Although this value for a column density is not uncommon for LMXB, the source needs to be in an hypothetically high state to be well detected by an Arcus-like instrument as the effective area of such an instrument would fall rapidly at the carbon edge.

Aluminum: the cosmic abundance of Al is significantly less than the main ID components (Fig. 1). However, Al in the ISM is almost completely depleted onto dust. We simulate here the contribution of Al\(_{2}\)O\(_3\) and MgAl\(_2\)O\(_4\) and compare them with a theoretical pure gas absorption (Fig. 3). At ISM temperature Al, if it was totally in gas form, would be distributed between Al\(_{1}\) and Al\(_{2}\). For this simulation we selected the bright LMXB \( (F_{2-10 \text{keV}} = 3 \times 10^{-9} \text{erg cm}^{-2} \text{s}^{-1}) \), GX 3+1 with its spectral parameters where obtained from Chandra-HETGS data (obsid 16492).

Sulfur: for the sulfur simulation, we selected GX 5-1, which is among the brightest LMXB in the Galaxy, with a \( F_{2-10 \text{keV}} = 2.5 \times 10^{-8} \text{erg cm}^{-2} \text{s}^{-1} \). The hydrogen column density is about \( 3.4 \times 10^{22} \) cm\(^{-2}\) (Zeegers et al. 2017). The depletion of sulfur is unknown in the diffuse ISM, but it has been estimated that could be up to \( \sim 46\% \) (Fig. 1 and Gry \\& Jenkins 2017). Here we simulate a more conservative 30\% depletion. Given the relatively low depletion of S, the XAFS features (Fig. 3) would be less evident in the data (Fig. 4).

Calcium: The X-ray spectrum will be sensitive to calcium extinction only if the intervening column density is sufficiently high. This is due to the relatively high energy position of the photoelectric edge. In Fig. 5, we simulated GX 340+00 \( (F_{2-10 \text{keV}} = 2.5 \times 10^{-8} \text{erg cm}^{-2} \text{s}^{-1}) \), with a column density of about \( N_H \sim 6.9 \times 10^{22} \) cm\(^{-2}\), obtained from HETG-Chandra data fitting (obsid 6632).

Titanium and nickel: for both these elements, we simulated an hypothetical source, for example near the GC, where also the occurrence of high column density molecular clouds is more frequent, that in outburst reaches a flux as high as GX 340+00 (Fig. 6). The column density must be sufficient to produce an edge-like modulation in the spectrum \( (N_H \sim 1.3 \times 10^{23} \) cm\(^{-2}\)).

4. Discussion

4.1. Carbon

In the simulation we included gas and the carbon forms that are believed to be most abundant (namely graphite, amorphous carbon and HAC). While the difference between graphite and amorphous carbon is subtle, HAC does have more distinctive features that may be more easily detected. The hydrogenation of carbon may point to either an environment protected from strong radiation or the presence of large grains, which are more resilient to radiation (Sect. 1). For illustrative purposes, we also include diamonds (orange dashed line in Fig. 2) in order to show the departure of this form of carbon from the shape of e.g. graphite. However, in practice, diamonds are believed to constitute no more than 5% of the carbon (Tielens 2001). Its realistic inclusion would be non-detectable (dashed-dotted blue line in figure). The same negligible effect may be produced by PAH, which we did not include in our simulation. The total amount along a line of sight is relatively low (Sect. 1) and the spectral features of PAH would be mixed with a more dominant amorphous carbon (or graphite) contribution. The sparse historical studies of PAH absorption profiles for the X-ray region have been recently revived (Reitsma et al. 2014, 2015). This will help in defining a shape for the summed contribution of the numerous different PAH in the ISM.

The carbon edge is not very sensitive to the size distribution of the grains (Draine 2003), therefore the MRN distribution may be adequate to describe the edge.

4.2. Aluminum and Calcium

The XIFU simulation shows that dust will be easily detectable, even if the edge itself produced a jump in the spectrum of only few %. However, from the Al edge alone, it would be difficult to distinguish among different compounds.

1. http://www.arcusray.org/
2. https://heasarc.gsfc.nasa.gov/docs/xrism/
3. http://www.the-athena-x-ray-observatory.eu/
We also note that contrary to other extinction profiles, the scattering peak, which appears as an emission-like feature before the edge-jump, is noticeable in Al. This peak is sensitive to the dust size distribution (Zeegers et al. 2017) and can be used, in principle, to estimate e.g. the mean grain size along the line of sight. As described above, we also tested the effect of the dust size distribution of Draine & Fraisse (2009) for Al$_2$O$_3$. The edge energy of Al lies in a zone sensitive to scattering (Draine 2003), therefore in Fig. 5 the large particles contribution is evident. As shown in Rogantini et al. (2018), the role of a substantial scattering contribution to the extinction not only forces the edge energy to shift, but may also modify the appearance of the edge absorption features. Grains containing seeds of Al and Ca, which are shielded from erosion in the ISM, are believed to be of large size, due to the several layers of coatings surrounding those seed elements (e.g. Clayton 1978 and Sect. 1.1). With future instruments we will therefore be able to test also the presence of larger particles for less abundant, but important, constituents of the ISM. The study of the Al edge will be however challenging, as Al is always a major component of X-ray space instruments (often in the form of foils). The extinction feature from Al in the ISM will be always blended with a relatively deep instrumental Al feature. This would need a careful calibration, adding uncertainty to the modeling.

Calcium is totally depleted in the ISM, therefore the main dust features will be detected (Fig. 5). However, calcium is mostly contained in silicates and aluminates, where oxygen is the main constituent. XAFS models shows that the first and main absorption feature is due to the nearest neighboring atom that the photoelectron wave will encounter (Lee & Ravel 2005). In the case simulated here, the absorption profile is dominated by oxygen (as is the case, to a lesser extend, in the Al edge), and only at higher energies are the secondary absorption features to be seen. For this reason the Ca inclusion in a specific silicate may be hard to disentangle through the observed spectrum. However, calcite (CaCO$_3$, dashed orange line), due to its different internal structure, will show a distinctive pattern, which may be in principle disentangled. This will help in determining whether this elusive compound (e.g. Kemper et al. 2002 and Sect. 1.1) may be present in the ISM. We tested the contribution of possible large grains on anorthite (blue dashed-dotted line in Fig. 5). The contribution of larger grains does not produce a well detectable feature.

4.3. Sulfur

In the diffuse ISM, sulfur is expected to have a modest depletion (Sect. 1.1). We use sulfur in conjunction with iron in the form of troilite, pyrrhotite and pyrite. FeS is a likely candidate for a diffuse interstellar environment, due to its inclusion in GEMS (Bradley 1994, Bradley et al. 1999). The line of sight towards GX 5-1, at distance of ∼9 kpc is likely to cross also molecular clouds and this would apply for any source located near the galactic center. The dust inclusion of sulfur in molecular clouds is still an open issue (Sect. 1.1). Some of the S must be associated to ices and carbon-hydrogen aggregates, while the rest may be in the form of FeS or atomic gas. Even the sum of all known S-bearing molecules would be unlikely to exceed few % of the total S abundance. Therefore any significant depletion detected by XRISM or XIFU would naturally point to the role of S in GEMS. We note that this amount of S depletion would still not procure visible deviations from the observed total dust spectral energy distribution (Köhler et al. 2014).

4.4. Titanium and Nickel

Due to its extremely low abundance in the Universe, titanium will be challenging to detect (Fig. 6). Nickel is about twenty times more abundant than Ti, however Ni will be also difficult to study (Fig. 7). The large column densities required to produce a Ni edge, will cause also strong absorption by iron, whose K absorption edge lies at 7100 eV, only 1.2 keV away from the one of nickel. Under the conditions of this simulation, the optical depth of iron will be around 18 times larger than the one of nickel. The net effect is that the Ni edge "sees" a continuum which is much lower than the one of the source, reducing the signal to noise ratio in that feature. Both titanium and nickel are however
completely depleted in most ISM environments, therefore even a column density estimate will be useful to constrain the abundance of these two elements, which are a product of explosions of both massive stars and white dwarfs.

5. Conclusion

In this paper we have shown how improved instrumental sensitivity and resolution will help in understanding new aspects of the composition of ID. Our results can be summarized as follows:

Future instruments, with characteristics similar to the Arcus mission, will be able to disentangle between the major components of carbon, namely amorphous carbon (or graphite) and hydrogenated carbon. The effect of minor constituents of C in the ISM (e.g. nano-diamonds and PAH) will be challenging to detect.

Instruments with improved capabilities at higher energies as Athena-XIFU or XRISM-Resolve, will be able to study absorption features that, due to their modest opacity, could not be investigated before. Simulations show that even a 1-6% jump in the transmission spectrum will be detected, allowing at least abundances measurements. For the low-cosmic abundance elements investigate in the E>1 keV band (namely Al, S and Ca), a full characterization (e.g. distinguishing among various silicate-like compounds) of the dust chemistry will be likely challenging. However, some main distinctions can be made:

– It will be possible to distinguish between calcium in carbonates and silicates around the Ca edge.
– For both Ca and Al the dust size distribution of these heavily depleted elements can be determined, with different precision depending on the instrument characteristics.
– It will be possible to determine the depletion of sulfur in the ISM. This in turn will help to clarify the S inclusion in GEMS, which are sometimes considered as one of the main forms of silicates in the ISM.

Finally, simulations show that Ti and Ni will be unaccessible to a detailed study even with next generation instruments considered here.

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Appendix A: Extinction profiles

We show here the extinction profiles in transmission, normalized for the continuum, of the compounds presented in this paper. Their formula and literature reference is reported in Table A.2. The instruments used for those measurements are reported in Table A.1.

Fig. A.1. The extinction profiles of the C, Al and S compounds presented in this paper. The dust column densities of the elements are the same used for the simulations. The edge energy is as reported in the literature as well as the original energy resolution.
Table A.1. Facilities and resolution of the literature laboratory measurements

| Ref                        | facility-beamline | element | resolution (eV) |
|----------------------------|-------------------|---------|-----------------|
| Albella et al. (1998)      | SSRL-8.2          | C       | 0.1             |
| Buijnsters et al. (2009)   | BessyII-SURICAT   | C       | ~0.05           |
| This work                  | Soleil-LUCIA      | Al      | 0.25            |
| Bonnin-Mosbah et al. (2002)| ESRF-ID21         | S       | 0.3             |
| Neuville et al. (2007)     | Soleil-LUCIA      | Ca      | 0.25            |
| Shin et al. (2013)         | PSLII-7D          | Ti      | ~0.09           |
| Van Loon et al. (2015)     | CLS-HXMA          | Ni      | 0.8             |
**Fig. A.2.** The extinction profiles of the Ca, Ti and Ni compounds presented in this paper. The dust column densities of the elements are the same used for the simulations. The edge energy is as reported in the literature as well as the original energy resolution.