Global Stability of Non-monotone Noncritical Traveling Waves for a Discrete Diffusion Equation with a Convolution Type Nonlinearity

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Abstract. This paper is concerned with the global stability of non-monotone traveling waves for a discrete diffusion equation with a monostable convolution type nonlinearity. It has been proved by Yang and Zhang (Sci. China Math. 61 (2018), 1789–1806) that all noncritical traveling waves (waves with speeds $c > c_*$, $c_*$ is minimal speed) are time-exponentially stable, when the initial perturbations around the waves are small. In this paper, we further prove that all traveling waves with large speed are globally stable, when the initial perturbations around the waves in a weighted Sobolev space can be arbitrarily large. The approaches adopted are the nonlinear Halanay’s inequality, the technical weighted energy method and Fourier’s transform.

1. Introduction

In this paper, we study the following spatially discrete diffusion equation with convolution type nonlinearity [8,28]:

$$
\frac{\partial u(t,x)}{\partial t} = D_2[u](t,x) - u(t,x) + \sum_{i \in \mathbb{Z}} K(i)g(u(t - \tau, x - i)), \quad t > 0, \quad x \in \mathbb{R}\n$$

with the initial data

$$
u(s, x) = u_0(s, x), \quad s \in [-\tau, 0], \quad x \in \mathbb{R},$$

where $\tau > 0$ and

$$D_2[u](t,x) = d[u(t, x + 1) - 2u(t,x) + u(t, x - 1)].$$

Equation (1.1) models the matured population dynamics of a single species with nonzero maturation delay. Here, $u(t,x)$ represents the mature population at the time $t$ and the location $x$, $d > 0$ is the coefficient of spacial diffusion, $\tau$ is the maturation delay,

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\[ \sum_{i \in \mathbb{Z}} K(i)g(u(t - \tau, x - i)) \] involves an infinite summation accounting for the non-local interaction, \( K \) is the non-negative weighted function satisfying

\[
K(i) = K(-i) \geq 0, \quad \sum_{i \in \mathbb{Z}} K(i) = 1 \quad \text{and} \quad \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda i} < \infty
\]

for any \( \lambda > 0 \), and \( g(\cdot) \) is the birth rate function satisfying

1. \( g(0) = 0, g(u^+) = u_+ \) for some positive constant \( u_+ \), \( g(u) > u \) for \( u \in (0, u_+) \), and \( g'(0) > 1 \) and \( g'(u_+) < 1 \);
2. \( g(u) > 0 \) has only one positive local maximum at the point \( u^* \in (0, u_+) \), and \( g(u) \) is increasing on \([0, u^*] \) and decreasing on \([u^*, +\infty) \);
3. \( g \in C^2[0, \infty) \) and \( |g'(u)| \leq g'(0) \) for \( u \in [0, \infty) \).

By (G1), we see that the equation (1.1) admits two constant equilibria \( u = u_- = 0 \) and \( u = u_+ \), where \( 0 \) is unstable and \( u_+ \) is stable. (G2) shows that \( g(u) \) is not monotone for \( u \in [0, u_+] \). Throughout this paper, we assume that

\[
\lim_{x \to \pm \infty} u_0(s, x) = u_\pm \quad \text{uniformly in} \quad s \in [-\tau, 0].
\]

A traveling wave of (1.1) is a special solution of the form \( u(t, x) = \phi(x + ct) \) connecting two equilibria \( u_- \) to \( u_+ \), where \( c > 0 \) is the wave speed. The wave profile equation of (1.1) is

\[
(1.4) \quad c\phi'(-\xi) - D_2[\phi](\xi) + \phi(\xi) = \sum_{i \in \mathbb{Z}} K(i)g(\phi(\xi - c\tau - i)), \quad \phi(-\infty) = 0, \quad \phi(+\infty) = u_+,
\]

where \( \xi = x + ct, \quad ' = d/d\xi, \quad D_2[\phi](\xi) = d[\phi(\xi + 1) - 2\phi(\xi) + \phi(\xi - 1)]. \) Moreover, if \( \phi(\xi) \) is monotone in \( \xi \in \mathbb{R} \), then it is called a traveling wavefront.

Equation (1.1) is a continuum version of the following lattice differential equation

\[
(1.5) \quad \frac{\partial u_n(t)}{\partial t} = d[u_{n+1}(t) - 2u_n(t) + u_{n-1}(t)] - u_n(t) + \sum_{i \in \mathbb{Z}} K(i)g(u_{n-i}(t - \tau)).
\]

Note that equations (1.1) and (1.5) possess the same wave profile equation (1.4). To the best of our knowledge, the existence of traveling waves and other properties, such as monotonicity and uniqueness of traveling waves of (1.5) were well studied. We refer the readers to [15, 16] for bistable case, [1, 7, 14, 25] for monotone monostable case, and [5, 6, 27, 30] for nonmonotone monostable case. However, little has been done for the stability of traveling waves of (1.1) and (1.5), when the function \( g \) is not monotone. More recently, Yang and Zhang [28] proved that all noncritical traveling waves of (1.1) (waves with speeds \( c > c_\ast, \ c_\ast \) is minimal speed) are time-exponentially stable, when the initial
perturbations around the waves are small. The method adopted in [28] is the technical weighted energy method [4, 11, 13, 19, 20, 26, 32]. We would like to mention that, Tian et al. [24] and Yang et al. [29], respectively, applied this method to prove the local stability of traveling waves of (1.1) when \( K(0) = 1 \) and \( K(i) = 0 \) for all \( i \neq 0 \). An interesting but also challenging question is whether these traveling waves of (1.1) with speed \( c \geq c^* \) are globally stable.

The global stability of traveling waves for various monostable evolution equations has been extensively investigated, see e.g., [2, 3, 9, 15, 17, 18, 23, 31], and references therein. There are main two methods: the squeezing technique [3, 15] for a class of discrete reaction-diffusion equations; and the weighted energy method together with the comparison principle developed by Mei and coauthors [17, 18] for the Nicholson’s blowflies equations, see also [9] for discrete reaction-diffusion equations with nonlocal delay effects, [23] for a discrete reaction-diffusion competition system. We should point out that the comparison principle is needed in above two methods. However, when the birth rate function \( g(u) \) in (1.1) is non-monotone, the comparison principle does not hold, and hence, the above two methods are not valid. In this paper, we shall apply the nonlinear Halanay’s inequality, the anti-weighted technique, Fourier’s transform and the boundedness estimate of traveling waves to establish the global stability of traveling waves. These methods were inspired by Mei et al. [4, 11, 22]. More precisely, the time-exponential decay of \( U(t, \xi) \) at \( \xi = +\infty \) of transformed equation can be obtained by using nonlinear Halanay’s inequality (see Proposition 3.3). To obtain the decay estimate of \( U(t, \xi) \) for \( \xi \in (-\infty, x_0] \) is a key step of this paper, where \( x_0 > 0 \) is large enough. We first take the anti-weighted technique to obtain a new equation (3.7). By some observation, the absolute value of the traveling waves of equation (3.7) can be bounded by the positive solution of a linear delayed discrete diffusion equation with constant coefficients, and Fourier’s transform can be applied to get the decay estimate for the solution of this linear delayed equation. Then we obtain the decay estimate of \( U(t, \xi) \) for \( \xi \in (-\infty, x_0] \). Combining above two decay estimates, we prove the global stability for the traveling waves with large speed.

The rest of this paper is organized as follows. In Section 2, we introduce some preliminaries and state our stability result. Section 3 is devoted to proving our main result on the global stability of traveling waves with large speed.

2. Preliminaries and the stability theorem

In this section, we first give the equivalent integral form of the initial value problem of (1.1), then recall the existence of traveling waves of (1.1), and finally state the main result on the stability of traveling waves.
First, we consider the initial value problem (1.1) with (1.2), i.e.,
\begin{equation}
\begin{cases}
\frac{\partial u(t,x)}{\partial t} = D_2[u](t,x) - u(t,x) + \sum_{i \in \mathbb{Z}} K(i) g(u(t-\tau,x-i)), \quad (t,x) \in \mathbb{R}_+ \times \mathbb{R}, \\
u(s,x) = u_0(s,x), \quad s \in [-\tau,0], \ x \in \mathbb{R}.
\end{cases}
\end{equation}

According to [10], with aid of the modified Bessel functions, the solution to the initial value problem
\begin{equation}
\begin{cases}
\frac{\partial u}{\partial t}(t,x) = d[u(t,x+1) - 2u(t,x) + u(t,x-1)], \quad (t,x) \in \mathbb{R}_+ \times \mathbb{R}, \\
u(0,x) = u_0(x), \quad x \in \mathbb{R}
\end{cases}
\end{equation}
can be expressed by
\begin{equation}
u(t,x) = (S(t)u_0)(x) = e^{-2dt} \sum_{m=-\infty}^{\infty} I_m(2dt)u_0(x-m),
\end{equation}
where $u_0(\cdot) \in L^\infty(\mathbb{R})$, $I_m$, $m \geq 0$ are defined as
\begin{equation}
I_m(t) = \sum_{k=0}^{\infty} \frac{(t/2)^{m+2k}}{k!(m+k)!},
\end{equation}
and $I_m(t) = I_{-m}(t)$ for $m < 0$. Moreover,
\begin{equation}
I_m'(t) = \frac{1}{2} [I_{m+1}(t) + I_{m-1}(t)], \quad \forall \ t \geq 0, \ m \in \mathbb{Z},
\end{equation}
and $I_m(0) = 0$ for $m \neq 0$ while $I_0(0) = 1$, and $I_m(t) \geq 0$ for any integer $m$ and $t \geq 0$. It is easy to see that
\begin{equation}
e^{-t} \sum_{m=-\infty}^{\infty} I_m(t) = e^{-t} [I_0(t) + 2I_1(t) + 2I_3(t) + \cdots] = 1.
\end{equation}
Clearly, (2.1) is equivalent to
\begin{equation}
u(t,x) = e^{-(2d+1)t} \sum_{m=-\infty}^{\infty} I_m(2dt)u_0(0,x-m)
+ \sum_{m=-\infty}^{\infty} \int_0^t e^{-(2d+1)(t-s)} I_m(2d(t-s)) \left( \sum_{i \in \mathbb{Z}} K(i) g(u(s-\tau,x-m-i)) \right) ds.
\end{equation}
In fact, by (G2), one has $g(u) \leq g(u_*)$ for $u \in [0, +\infty)$. Then by [10 Lemma 2.1], we can
differentiate the series on $t$ variable in (2.4). Then we obtain
\[
\frac{\partial u(t, x)}{\partial t} = -(2d + 1) e^{-(2d+1)t} \sum_{m=-\infty}^{\infty} I_m(2dt)u_0(0, x - m) \\
+ e^{-(2d+1)t} \sum_{m=-\infty}^{\infty} 2dI'_m(2dt)u_0(0, x - m) \\
+ \sum_{m=-\infty}^{\infty} I_m(0) \sum_{i \in \mathbb{Z}} K(i)g(u(t - \tau, x - m - i)) \\
- (2d + 1) \sum_{m=-\infty}^{\infty} \int_0^t e^{-(2d+1)(t-s)} I_m(2dt - s) \left( \sum_{i \in \mathbb{Z}} K(i)g(u(s - \tau, x - m - i)) \right) ds \\
+ \sum_{m=-\infty}^{\infty} \int_0^t e^{-(2d+1)(t-s)} 2dI'_m(2dt - s) \left( \sum_{i \in \mathbb{Z}} K(i)g(u(s - \tau, x - m - i)) \right) ds \\
= d[u(t, x + 1) - 2u(t, x) + u(t, x - 1)] - u(t, x) + \sum_{i \in \mathbb{Z}} K(i)g(u(t - \tau, x - i)),
\]
where we have used the recurrence relation (2.2).

The characteristic function for (1.4) with respect to the trivial equilibrium 0 can be represented by
\[
P(c, \lambda) = c\lambda - d(e^\lambda + e^{-\lambda} - 2) + 1 - g'(0)G(\lambda),
\]
where
\[
G(\lambda) = \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} < \infty.
\]
One can easily show that the following result holds.

**Lemma 2.1.** Assume that $g'(0) > 1$. Then there exist $\lambda_* > 0$ and $c_* > 0$ such that
\[
P(c_*, \lambda_*) = 0 \quad \text{and} \quad \frac{\partial}{\partial \lambda} P(c_*, \lambda)\bigg|_{\lambda=\lambda_*} = 0.
\]
Furthermore, if $c > c_*$, then $P(c, \lambda) = 0$ has two distinct positive real roots $\lambda_1(c)$ and $\lambda_2(c)$ with $\lambda_1(c) < \lambda_* < \lambda_2(c)$, and $P(c, \lambda) > 0$ for $\lambda \in (\lambda_1(c), \lambda_2(c))$.

The existence of traveling waves has been obtained by Yang and Zhang, see [28, Theorem 2.3].

**Proposition 2.2.** Assume that (G1)–(G3) hold. Then for every $c > c_*$, (1.1) admits a traveling wave $u(t, x) = \phi(x + ct)$ satisfying $\phi(-\infty) = 0$ and $u^*_+ < \liminf_{\xi \to +\infty} \phi(\xi) \leq \limsup_{\xi \to +\infty} \phi(\xi) \leq u^*_+ > u_+$ for some positive constant $u^*_+$. 
Remark 2.3. We should point out if we further assume that the following assumption (G4) holds, then \( \phi(+\infty) = u_+ \).

\[
\text{G4)} \quad g(u) < 2u_+ - u \text{ for } u \in [u_-^*, u_+] \text{ and } g(u) > 2u_+ - u \text{ for } u \in (u_+, u_+^*], \text{ where } 0 < u_-^* \leq u_+ \leq u_+^* \text{ and } g_+(u_+) = du_+^*, \text{ and } g_-(u_-) = du_-^*, \text{ where } g_+(u) := \min_{w \in [0, u]} g(w) \text{ and } g_-(u) := \min_{w \in [u, u_+^*]} g(w).
\]

Notations. \( C > 0 \) denotes a generic constant, while \( C_i \) \( (i = 1, 2, \ldots) \) represents a specific constant. Let \( I \) be an interval, typically \( I = \mathbb{R} \). Denote by \( L^1(I) \) the space of integrable functions defined on \( I \), and \( W^{k,1}(I) \) \( (k \geq 0) \) the Sobolev space of the \( L^1 \)-functions \( f(x) \) defined on the interval \( I \) whose derivatives \( \frac{d^n}{dx^n} f \) \( (n = 1, \ldots, k) \) also belong to \( L^1(I) \). Let \( L^1_w(I) \) be the weighted \( L^1 \)-space with a weight function \( w(x) > 0 \) and its norm is defined by

\[
\| f \|_{L^1_w(I)} = \int_I w(x)|f(x)| \, dx,
\]

\( W^{k,1}_w(I) \) be the weighted Sobolev space with the norm given by

\[
\| f \|_{W^{k,1}_w(I)} = \sum_{i=0}^k \int_I w(x) \left| \frac{d^i f(x)}{dx^i} \right| \, dx.
\]

Let \( T > 0 \) be a number and \( B \) be a Banach space. We denote by \( C([0, T]; B) \) the space of the \( B \)-valued continuous functions on \([0, T]\), and by \( L^1([0, T]; B) \) the space of the \( B \)-valued \( L^1 \)-functions on \([0, T]\). The corresponding spaces of the \( B \)-valued functions on \([0, \infty)\) are defined similarly. For any function \( f(x) \), its Fourier transform is defined by

\[
\mathcal{F}[f](\eta) = \hat{f}(\eta) = \int_{\mathbb{R}} e^{-i\xi \eta} f(x) \, dx
\]

and the inverse Fourier transform is given by

\[
\mathcal{F}^{-1}[\hat{f}](x) = \frac{1}{2\pi} \int_{\mathbb{R}} e^{i\xi \eta} \hat{f}(\eta) \, d\eta,
\]

where \( i \) is the imaginary unit, \( i^2 = -1 \).

Define a weight function

\[
\omega(\xi) := e^{-2\lambda \xi}, \quad \xi \in \mathbb{R},
\]

where \( \lambda = c_* \). Notice that \( \lim_{\xi \to -\infty} \omega(\xi) = +\infty \) and \( \lim_{\xi \to +\infty} \omega(\xi) = 0 \), since \( \lambda = c_* > 0 \).

For the sake of simplicity, we denote

\[
L_1 := \frac{d}{2}(e^{c_*} + e^{-c_*} - 2), \quad L_2 := \frac{e^{d(e^{c_*} + e^{-c_*})}}{2} \sum_{i \in \mathbb{Z}} K(i) e^{-c_* i}.
\]

Now we state the stability result of traveling waves of (1.1) with a general nonmonotone function \( g(u) \).
Theorem 2.4 (Global stability). Assume that (G1)–(G3) hold. For any given traveling wave \( \phi(x + ct) = \phi(\xi) \) connecting 0 and \( u_+ \) with

\[
c > \tilde{c} := \max \left\{ c_*, \frac{L_1 - 1 + L_2}{c_*} \right\},
\]

whether it is monotone or non-monotone, when the initial perturbation satisfies \( u_0 - \phi \in \mathcal{C}_{\text{unif}}[-\tau, 0] \cap C([\tau, 0]; W^{1,1}_\omega(\mathbb{R})) \) and \( \partial_\xi (u_0 - \phi) \in L^1([\tau, 0]; L^1_\omega(\mathbb{R})) \), we have the following global stability

\[
\sup_{x \in \mathbb{R}} |u(t, x) - \phi(x + ct)| \leq Ce^{-\mu t},
\]

where \( \mu \) and \( C \) are positive numbers, and \( \mathcal{C}_{\text{unif}}[-\tau, T] \) is the uniformly continuous space, for \( 0 < T \leq \infty \), defined by

\[
\mathcal{C}_{\text{unif}}[-\tau, T] = \left\{ u \in C(\mathbb{T} \times \mathbb{R}) \text{ such that } \lim_{x \to +\infty} u(t, x) \text{ exists uniformly in } t \in [-\tau, T] \right\}.
\]

3. The global stability of traveling waves

In this section, we are devoted to the proof of global stability of those monotone or non-monotone traveling waves of (1.1) with speed \( c > \tilde{c} \), when \( g \) is non-monotone.

Let \( \phi(x + ct) = \phi(\xi) \) be any given traveling wave with speed \( c > c_* \), and define

\[
U(t, \xi) := u(t, x) - \phi(x + ct) = u(t, \xi - ct) - \phi(\xi),
\]

\[
U_0(s, \xi) := u_0(s, x) - \phi(x + cs).
\]

Then it follows from (1.1) and (1.4) that \( U(t, \xi) \) satisfies

\[
\frac{\partial U}{\partial t} + c \frac{\partial U}{\partial \xi} - \mathcal{D}_2[U] + U = \sum_{i \in \mathbb{Z}} K(i) Q(U(t - \tau, \xi - ct + i)), \quad (t, \xi) \in \mathbb{R}_+ \times \mathbb{R},
\]

\[
U(s, \xi) = U_0(s, \xi), \quad s \in [\tau, 0], \quad \xi \in \mathbb{R},
\]

where

\[
Q(U) := g(\phi + U) - g(\phi) = g'('\tilde{\phi}')U
\]

for some \( \tilde{\phi} \) between \( \phi \) and \( \phi + U \), with \( \phi = \phi(\xi - ct + i) \) and \( U = U(t - \tau, \xi - ct - i) \) for \( c > c_* \).

We first prove the existence and uniqueness of solution to the initial value problem (3.1) in the uniformly continuous space \( \mathcal{C}_{\text{unif}}[-\tau, \infty) \).

Lemma 3.1. Assume that (G1)–(G3) hold. If the initial perturbation \( U_0 \in \mathcal{C}_{\text{unif}}[-\tau, 0] \) for \( c > c_* \), then the solution \( U(t, \xi) \) of perturbed equation (3.1) is unique and time-globally exists in \( \mathcal{C}_{\text{unif}}[-\tau, \infty) \).
Proof. Let \( V(t, x) = u(t, x) - \phi(x + ct) \). It is clear that \( V(t, x) = U(t, \xi) \) and satisfies

\[
\begin{aligned}
\frac{\partial V}{\partial t} - D_2[V] + V = \sum_{i \in \mathbb{Z}} K(i) Q(V(t - \tau, x - i)), \quad (t, x) \in \mathbb{R}_+ \times \mathbb{R}, \\
V(s, x) = u_0(s, x) - \phi(x + cs) := V_0(s, x), \quad (s, x) \in [-\tau, 0] \times \mathbb{R}.
\end{aligned}
\] (3.3)

Thus, the global existence and uniqueness of solutions of (3.1) are transformed into that of (3.3).

When \( t \in [0, \tau] \), we have \( t - \tau \in [-\tau, 0] \) and \( V(t - \tau, x - i) = V_0(t - \tau, x - i) \). Consequently, the solution of (3.3) can be explicitly and uniquely solved by

\[
V(t, x) = e^{-(2d+1)t} \sum_{m=-\infty}^{\infty} I_m(2dt) V_0(0, x - m)
+ \sum_{m=-\infty}^{\infty} \int_{0}^{t} e^{-(2d+1)(t-s)} I_m(2d(t-s)) \left( \sum_{i \in \mathbb{Z}} K(i) Q(V_0(s - \tau, x - m - i)) \right) ds
\] (3.4)

for \( t \in [0, \tau] \).

Since \( U_0 \in C_{\text{unif}}[-\tau, 0] \), \( \lim_{\xi \to +\infty} U_0(\xi, t) \) exists uniformly in \( t \in [-\tau, 0] \), which implies \( \lim_{x \to +\infty} V_0(t, x) \) exists uniformly in \( t \in [-\tau, 0] \). Denote \( V_0(t, \infty) = \lim_{x \to +\infty} V_0(t, x) \). Taking the limit \( x \to +\infty \) to (3.4) yields

\[
\lim_{x \to +\infty} V(t, x) = e^{-(2d+1)t} \sum_{m=-\infty}^{\infty} I_m(2dt) \lim_{x \to +\infty} V_0(0, x - m)
+ \sum_{m=-\infty}^{\infty} \int_{0}^{t} e^{-(2d+1)(t-s)} I_m(2d(t-s)) \left( \sum_{i \in \mathbb{Z}} K(i) \lim_{x \to +\infty} Q(V_0(s - \tau, x - m - i)) \right) ds
= e^{-t} V_0(0, \infty) + \int_{0}^{t} e^{-(t-s)} Q(V_0(s - \tau, \infty)) \sum_{m=-\infty}^{\infty} e^{-2d(t-s)} I_m(2d(t-s)) ds
=: \mathcal{V}_1(t) \quad \text{uniformly in } t \in [0, \tau].
\] (3.5)

Here we use the fact that \( \sum_{i \in \mathbb{Z}} K(i) = 1 \) and (2.3). Thus, we obtain that \( V \in C_{\text{unif}}[-\tau, \tau] \), i.e., \( U \in C_{\text{unif}}[-\tau, \tau] \).

When \( t \in [\tau, 2\tau] \), equation (3.3) with the initial data \( V(s, x) \) for \( s \in [0, \tau] \) is still linear because the source term \( \sum_{i \in \mathbb{Z}} K(i) Q(V(t - \tau, x - i)) \) is known due to \( t - \tau \in [0, \tau] \) and \( V(t - \tau, x) \) is solved in (3.4). Hence, the solution \( V(t, x) \) for \( t \in [\tau, 2\tau] \) is uniquely and
explicitly given by
\[ V(t, x) = e^{-(2d+1)(t-\tau)} \sum_{m=-\infty}^{\infty} \mathcal{I}_m(2d(t-\tau)) V(\tau, x-m) \]
\[ + \sum_{m=-\infty}^{\infty} \int_{\tau}^{t} e^{-(2d+1)(t-s)} \mathcal{I}_m(2d(t-s)) \left( \sum_{i \in \mathbb{Z}} K(i) Q(V(s-\tau, x-m-i)) \right) \, ds. \]

Similarly, by (3.5), we have
\[ \lim_{x \to +\infty} V(t, x) = e^{-(2d+1)(t-\tau)} \sum_{m=-\infty}^{\infty} \mathcal{I}_m(2d(t-\tau)) \lim_{x \to +\infty} V(\tau, x-m) \]
\[ + \sum_{m=-\infty}^{\infty} \int_{\tau}^{t} e^{-(2d+1)(t-s)} \mathcal{I}_m(2d(t-s)) \left( \sum_{i \in \mathbb{Z}} K(i) \lim_{x \to +\infty} Q(V(s-\tau, x-m-i)) \right) \, ds \]
\[ = e^{-(t-\tau)} \mathcal{V}_1(\tau) + \int_{\tau}^{t} e^{-(t-s)} Q(\mathcal{V}_1(s-\tau)) \sum_{m=-\infty}^{\infty} e^{-2d(t-s)} \mathcal{I}_m(2d(t-s)) \, ds \]
\[ =: \mathcal{V}_2(t) \text{ uniformly in } t \in [\tau, 2\tau]. \]

Repeating the above procedure for \( t \in [n\tau, (n+1)\tau] \) with \( n = 2, 3, \ldots \), we can obtain that \( U \in \mathcal{C}_{\text{unif}}[-\tau, (n+1)\tau] \) uniquely exists. Step by step, we can finally prove the global existence and uniqueness of solution \( U \in \mathcal{C}_{\text{unif}}[-\tau, +\infty) \).

Now we state the stability result for the perturbed equation (3.1), which automatically implies Theorem 2.4.

**Proposition 3.2.** Assume that (G1)–(G3) hold. If the initial perturbation 
\[ U_0 \in \mathcal{C}_{\text{unif}}[-\tau, 0] \cap C([-\tau, 0]; W^{1,1}_\omega(\mathbb{R})), \quad \partial_x U_0 \in L^1([-\tau, 0]; L^1_\omega(\mathbb{R})), \]
then when \( c > \tilde{c} \), it holds
\[ \sup_{\xi \in \mathbb{R}} |U(t, \xi)| \leq Ce^{-\mu t}, \quad t > 0 \]
for some positive constant \( \mu \).

Proposition 3.2 is a straightforward consequence of the following Propositions 3.3 and 3.11. In the rest of this section, we always assume the assumptions (G1)–(G3) hold.

**Proposition 3.3.** Assume that \( U_0 \in \mathcal{C}_{\text{unif}}[-\tau, 0] \). Then there exists a large number \( x_0 \gg 1 \) such that the solution \( U(t, \xi) \) of (3.1) satisfies
\[ \sup_{\xi \in [x_0, \infty)} |U(t, \xi)| \leq Ce^{-\mu_1 t}, \quad t > 0 \]
for some \( 0 < \mu_1 < 1 \).
Proof. Since \( U_0 \in C_{\text{unif}}[-\tau,0] \), by Lemma 3.1, we have \( U \in C_{\text{unif}}[-\tau,\infty) \), which implies \( \lim_{\xi \to +\infty} U(t,\xi) = U(t,\infty) =: z(t) \) exists uniformly for \( t \in [-\tau,\infty) \). Taking the limit \( \xi \to \infty \) to (3.1), one has
\[
\frac{d}{dt}z(t) + z(t) - g'(u_+)z(t - \tau) = P(z(t - \tau)), \quad z(s) = z_0(s), \quad s \in [-\tau,0],
\]
where
\[
P(z) = g(u_+ + z) - g(u_+) - g'(u_+)z.
\]
It follows from the Taylor expansion that \(|P(z)| \leq C|z|^2\) for some positive constant \( C \). By the nonlinear Halanay’s inequality given in [13], we obtain
\[
|U(t,\infty)| = |z(t)| \leq Ce^{-\mu_1 t}, \quad t > 0
\]
for some constant \( 0 < \mu_1 < 1 \), provided \( \|z_0\|_{L^\infty} \ll 1 \).

By the continuity and the uniform convergence of \( U(t,\xi) \) as \( \xi \to +\infty \), there exists a large number \( x_0 \gg 1 \) such that (3.6) implies
\[
\sup_{\xi \in [x_0,\infty)} |U(t,\xi)| \leq Ce^{-\mu_1 t}, \quad t > 0,
\]
provided \( \sup_{\xi \in [x_0,\infty)} |U_0(s,\xi)| \ll 1 \) for \( s \in [-\tau,0] \). Such a smallness for the initial perturbation \( U_0 \) near \( \xi \to +\infty \) can be automatically verified, since \( \lim_{x \to +\infty} u_0(s,x) = u_+ \) (see (1.3)), which implies \( \lim_{\xi \to +\infty} U_0(s,\xi) = \lim_{\xi \to +\infty}[u_0(s,\xi) - \phi(\xi)] = u_+ - u_+ = 0 \) uniformly for \( s \in [-\tau,0] \).

It remains to establish a priori decay estimate of \( \sup_{\xi \in (-\infty,x_0]} |U(t,\xi)| \). We shall use the anti-weighted technique [4,11] together with Fourier’s transform to treat this problem. Define
\[
\tilde{U}(t,\xi) = \sqrt{\omega(\xi)}U(t,\xi + x_0) = e^{-\lambda\xi}U(t,\xi + x_0),
\]
where \( \lambda = c_* \). It is easy to see that \( \tilde{U}(t,\xi) \) satisfies
\[
\begin{cases}
\frac{\partial \tilde{U}}{\partial t} + c_1 \frac{\partial \tilde{U}}{\partial \xi} + c_1 \tilde{U}(t,\xi) - d\lambda \tilde{U}(t,\xi + 1) - d\lambda \tilde{U}(t,\xi - 1) \\
\sum_{i \in \mathbb{Z}} K(i) \tilde{Q}(\tilde{U}(t - \tau,\xi - c\tau - i)),
\end{cases}
\]
(3.7)
\[
\tilde{U}(s,\xi) = \sqrt{\omega(\xi)}U(s,\xi + x_0) =: \tilde{U}_0(s,\xi), \quad s \in [-\tau,0], \quad \xi \in \mathbb{R},
\]
where \( c_1 = c\lambda + 2d + 1 \) and
\[
\tilde{Q}(\tilde{U}) = e^{-\lambda\xi}Q(U).
\]
By (3.2), we obtain
\[
\tilde{Q}(\tilde{U}(t - \tau,\xi - c\tau - i)) = e^{-\lambda\xi}Q(U(t - \tau,\xi + x_0 - c\tau - i))
\]
(3.8)
\[
= e^{-\lambda\xi}g'(\phi)U(t - \tau,\xi + x_0 - c\tau - i)
\]
\[
= e^{-\lambda(i + c\tau)}g'(\phi)\tilde{U}(t - \tau,\xi - c\tau - i)
\]
for some function $\tilde{\phi}$ between $\phi$ and $\phi + U$. By (G3), we further have
$$|\tilde{Q}(\tilde{U}(t-\tau, \xi - ct - i))| \leq g'(0)e^{-\lambda(i+ct)}|\tilde{U}(t-\tau, \xi - ct - i)|.$$ 

Since $g'(s)$ can be negative for $s \in (0, u_+)$, then the solution $\tilde{U}$ of (3.7) with the non-linear term (3.8) may be oscillation around $u_+$ when the time delay $\tau$ is large. At the same time, the comparison principle may not hold for (3.7). Hence, the monotonic technique cannot be applied. Note that the coefficient $g'(\tilde{\phi})$ in (3.8) is variable. As such, we are unable to derive the decay estimate directly by applying Fourier’s transform. In this paper, we shall provide some new idea to establish the decay estimate of $\sup_{\xi \in [-\infty, x_0]} |U(\xi, t)|$. By replacing $g'(\tilde{\phi})$ in (3.7) with $g'(0)$, we obtain a linear delayed discrete diffusion equation (3.9)

$$\frac{\partial U^+}{\partial t} + c\frac{\partial U^+}{\partial \xi} + c_1 U^+(t, \xi) - de^\lambda U^+(t, \xi + 1) - de^{-\lambda} U^+(t, \xi - 1)$$
$$= g'(0) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+ct)}U^+(t-\tau, \xi - ct - i), \quad (t, \xi) \in \mathbb{R}_+ \times \mathbb{R},$$

$$U^+(s, \xi) = U^+_0(s, \xi) \geq 0, \quad s \in [-\tau, 0], \ \xi \in \mathbb{R}.$$

We first prove that the solution $\tilde{U}$ of (3.7) can be bounded by the solution $U^+$ of (3.9), i.e., $|\tilde{U}(t, \xi)| \leq U^+(t, \xi)$ for $(t, \xi) \in \mathbb{R}_+ \times \mathbb{R}$. Since (3.9) is a linear equation, we can apply Fourier’s transform to investigate the decay estimate of $U^+$. Then the decay estimate of $\tilde{U}$ can be obtained.

In order to obtain the crucial boundedness estimate of $\tilde{U}$, we need the following maximum principle.

**Lemma 3.4.** Let $T > 0$. For any $a_1, a_2 \in \mathbb{R}$ and $\nu > 0$, if the bounded function $v$ satisfies

$$\frac{\partial v}{\partial t} + a_1 \frac{\partial v}{\partial \xi} + a_2 v - de^\nu v(t, \xi + 1) - de^{-\nu} v(t, \xi - 1) \geq 0, \quad (t, \xi) \in (0, T] \times \mathbb{R},$$

then $v(t, \xi) \geq 0$ for all $(t, \xi) \in (0, T] \times \mathbb{R}$.

**Proof.** Let $\tilde{v}(t, \xi) = v(t, \xi + a_1 t)$. Then by (3.10), $\tilde{v}(\xi, t)$ satisfies

$$\frac{\partial \tilde{v}}{\partial t} + a_2 \tilde{v} - de^\nu \tilde{v}(t, \xi + 1) - de^{-\nu} \tilde{v}(t, \xi - 1) \geq 0, \quad (t, \xi) \in (0, T] \times \mathbb{R},$$

$$\tilde{v}(0, \xi) = v(0, \xi) \geq 0, \quad \xi \in \mathbb{R}.$$

Let $w(t, \xi) = e^{kt}\tilde{v}(t, \xi)$, where $k > 0$ is chosen such that $p_0 := k - a_2 > 0$. Then we have

$$\frac{\partial w(t, \xi)}{\partial t} \geq de^\nu \tilde{w}(t, \xi + 1) + de^{-\nu} \tilde{w}(t, \xi - 1) + p_0 w(t, \xi), \quad (t, \xi) \in (0, T] \times \mathbb{R}.$$

Denote $0 < T_0 < \min \left\{ T, \frac{1}{2(d e^\nu + d e^{-\nu} + p_0)} \right\}$. We now prove that $w \geq 0$ in $(0, T_0) \times \mathbb{R}$. Suppose on the contrary that there are $\tilde{t} \in (0, T_0)$ and $\tilde{\xi} \in \mathbb{R}$ such that $w(\tilde{t}, \tilde{\xi}) < 0$. Then we have
Then it follows from (3.9) that
\[ w(t_n, \xi_n) - w(0, \xi_n) \geq d \int_0^{t_n} \left( e^{\nu} w(s, \xi_n + 1) + e^{-\nu} w(s, \xi_n - 1) \right) ds + \int_0^{t_n} p_0 w(s, \xi_n) ds \]
\[ \geq (de^\nu + de^{-\nu} + p_0) w_{\text{inf}} t_n \]
\[ \geq (de^\nu + de^{-\nu} + p_0) w_{\text{inf}} t^*. \]

Note that \( w(0, \xi_n) \geq 0 \) for \( n = 1, 2, \ldots \). Then we have \( w(t_n, \xi_n) \geq (de^\nu + de^{-\nu} + p_0) w_{\text{inf}} t^* \).

Letting \( n \to \infty \), we obtain
\[ w_{\text{inf}} \geq (de^\nu + de^{-\nu} + p_0) w_{\text{inf}} t^* > \frac{1}{2} w_{\text{inf}}, \]
a contradiction. Hence, \( v(t, \xi) \geq 0 \) for all \((t, \xi) \in (0, T] \times \mathbb{R}\). \( \square \)

**Lemma 3.5.** If \( U_0^+(s, \xi) \geq 0 \) for \((s, \xi) \in [-\tau, 0] \times \mathbb{R}\), then \( U^+(t, \xi) \geq 0 \) for \((t, \xi) \in [-\tau, +\infty) \times \mathbb{R}\).

**Proof.** For \( t \in [0, \tau] \), we have \( t - \tau \in [-\tau, 0] \) and
\[
g'(0) \sum_{i \in \mathbb{Z}} K(i) e^{-\lambda(i+c\tau)} U^+(t - \tau, \xi - c\tau - i) = g'(0) \sum_{i \in \mathbb{Z}} K(i) e^{-\lambda(i+c\tau)} U_0^+(t - \tau, \xi - c\tau - i) \geq 0.
\]

Then it follows from (3.9) that
\[
\frac{\partial U^+}{\partial t} + c \frac{\partial U^+}{\partial \xi} + c_1 U^+(t, \xi) - de^{\lambda} U^+(t, \xi + 1) - de^{-\lambda} U^+(t, \xi - 1) \geq 0,
\]
which implies \( U^+(t, \xi) \geq 0 \) for \( t \in [0, \tau] \) due to Lemma 3.4. Repeating this procedure step by step, we can prove that \( U^+(t, \xi) \geq 0 \) for \( t \in [n\tau, (n+1)\tau] \). Furthermore, we have \( U^+(t, \xi) \geq 0 \) for \((t, \xi) \in \mathbb{R}_+ \times \mathbb{R}\). \( \square \)

The following result shows the boundedness estimate for the solution \( \tilde{U}(t, \xi) \) of (3.7).

**Lemma 3.6.** Let \( \tilde{U}(t, \xi) \) and \( U^+(t, \xi) \) be the solutions of (3.7) and (3.9), respectively. When
\[
|\tilde{U}_0(s, \xi)| \leq U_0^+(s, \xi), \quad (s, \xi) \in [-\tau, 0] \times \mathbb{R},
\]
then
\[
|\tilde{U}(t, \xi)| \leq U^+(t, \xi), \quad (t, \xi) \in \mathbb{R}_+ \times \mathbb{R}.
\]
Proof. We first prove $|\tilde{U}(t, \xi)| \leq U^+(t, \xi)$ for $(t, \xi) \in [0, \tau] \times \mathbb{R}$. When $t \in [0, \tau]$, namely, $t - \tau \in [-\tau, 0]$, by (3.12), we have

\begin{equation}
|\tilde{U}(t - \tau, \xi - c\tau - i)| = |\tilde{U}_0(t - \tau, \xi - c\tau - i)| \\
\leq U^+_0(t - \tau, \xi - c\tau - i) = U^+(t - \tau, \xi - c\tau - i).
\end{equation}

Let

$$V^-(t, \xi) := U^+(t, \xi) - \tilde{U}(t, \xi) \quad \text{and} \quad V^+(t, \xi) := U^+(t, \xi) + \tilde{U}(t, \xi).$$

By (3.7), (3.8) and (3.9), we see that $V^-(t, \xi)$ satisfies

\begin{equation}
\frac{\partial V^-}{\partial t} + c \frac{\partial V^-}{\partial \xi} + c_1 V^-(t, \xi) - de^\lambda V^-(t, \xi + 1) - de^{-\lambda} V^-(t, \xi - 1) \\
= g'(0) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} U^+(t - \tau, \xi - c\tau - i) \\
- g'(\bar{\phi}) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} \tilde{U}(t - \tau, \xi - c\tau - i) \\
\geq g'(0) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} U^+(t - \tau, \xi - c\tau - i) \\
- \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} |g'(\bar{\phi})||\tilde{U}(t - \tau, \xi - c\tau - i)|.
\end{equation}

Furthermore, by (3.13) and (G3), we obtain from (3.14) that

\begin{equation}
\frac{\partial V^-}{\partial t} + c \frac{\partial V^-}{\partial \xi} + c_1 V^-(t, \xi) - de^\lambda V^-(t, \xi + 1) - de^{-\lambda} V^-(t, \xi - 1) \geq 0
\end{equation}

for $(t, \xi) \in [0, \tau] \times \mathbb{R}$. Thus, (3.15) with the initial data $V_0^-(s, \xi) = U^+_0(s, \xi) - \tilde{U}_0(s, \xi) \geq 0$ reduces to

$$\frac{\partial V^-}{\partial t} + c \frac{\partial V^-}{\partial \xi} + c_1 V^-(t, \xi) - de^\lambda V^-(t, \xi + 1) - de^{-\lambda} V^-(t, \xi - 1) \geq 0, \quad V_0^-(0, \xi) \geq 0.$$

By Lemma 3.4, we have

\begin{equation}
V^-(t, \xi) = U^+(t, \xi) - \tilde{U}(t, \xi) \geq 0 \quad \text{for} \quad (t, \xi) \in [0, \tau] \times \mathbb{R}.
\end{equation}

On the other hand, $V^+(t, \xi)$ satisfies

\begin{equation}
\frac{\partial V^+}{\partial t} + c \frac{\partial V^+}{\partial \xi} + c_1 V^+(t, \xi) - de^\lambda V^+(t, \xi + 1) - de^{-\lambda} V^+(t, \xi - 1) \\
= g'(0) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} U^+(t - \tau, \xi - c\tau - i) \\
- g'(\bar{\phi}) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} \tilde{U}(t - \tau, \xi - c\tau - i)
\end{equation}
\[
\geq g'(0) \sum_{i \in \mathbb{Z}} K(i) e^{-\lambda(i+ct)} U^+(t - \tau, \xi - c\tau - i) \\
- \sum_{i \in \mathbb{Z}} K(i) e^{-\lambda(i+ct)} |g'(\tilde{\phi})| \tilde{U}(t - \tau, \xi - c\tau - i) \\
\geq 0, \quad t \in [0, \tau].
\]

Similarly, we can prove that (3.17) with the initial data \( V^+_0(s, \xi) = U^+_0(s, \xi) + \tilde{U}_0(s, \xi) \geq 0 \) implies

\[
V^+(t, \xi) = U^+(t, \xi) + \tilde{U}(t, \xi) \geq 0 \quad \text{for} \quad (t, \xi) \in [0, \tau] \times \mathbb{R}.
\]

Combining (3.16) and (3.18), we have

\[
\tilde{U}(t, \xi) \leq U^+(t, \xi) \quad \text{for} \quad (t, \xi) \in [0, \tau] \times \mathbb{R}.
\]

Next, when \( t \in [\tau, 2\tau] \), namely, \( t - \tau \in [0, \tau] \), based on (3.19), we can similarly prove

\[
V^-(t, \xi) = U^+(t, \xi) - \tilde{U}(t, \xi) \geq 0,
\]

\[
V^+(t, \xi) = U^+(t, \xi) + \tilde{U}(t, \xi) \geq 0
\]

for \( (t, \xi) \in [\tau, 2\tau] \times \mathbb{R} \), namely,

\[
\tilde{U}(t, \xi) \leq U^+(t, \xi) \quad \text{for} \quad (t, \xi) \in [\tau, 2\tau] \times \mathbb{R}.
\]

Repeating this procedure, we then further prove

\[
\tilde{U}(t, \xi) \leq U^+(t, \xi) \quad \text{for} \quad (t, \xi) \in [n\tau, (n+1)\tau] \times \mathbb{R}, \ n = 1, 2, \ldots,
\]

which implies

\[
\tilde{U}(t, \xi) \leq U^+(t, \xi) \quad \text{for} \quad (t, \xi) \in \mathbb{R}_+ \times \mathbb{R}.
\]

This completes the proof. \( \square \)

Now we are in a position to establish the decay estimate of the solution \( U^+ \) of (3.9). We first recall some properties of the solutions to the delayed ODE.

Lemma 3.7. [12, Theorem 1] Let \( z(t) \) be the solution to the following scalar differential equation with delay

\[
\begin{aligned}
\frac{d}{dt} z(t) + k_1 z(t) &= k_2 z(t - \tau), \quad t \geq 0, \ \tau > 0, \\
z(s) &= z_0(s), \quad s \in [-\tau, 0].
\end{aligned}
\]

Then

\[
z(t) = e^{-k_1(t+\tau)} e^{\int_{-\tau}^t z_0(-s) ds} + \int_{-\tau}^t e^{-k_1(t-s)} e^{\int_{-\tau}^s z_0(-r) ds} [z_0(s) + k_1 z_0(s)] ds,
\]
where
\[ \bar{k}_2 = k_2 e^{k_1 \tau} \]
and \( e_{\tau}^{\bar{k}_2 t} \) is the so-called delayed exponential function in the form
\[
e_{\tau}^{\bar{k}_2 t} = \begin{cases} 
0, & -\infty < t < -\tau, \\
1, & -\tau \leq t < 0, \\
1 + \bar{k}_2 \frac{t}{\tau}, & 0 \leq t < \tau, \\
1 + \bar{k}_2 \frac{t}{\tau} + \bar{k}_2^2 \frac{(t-\tau)^2}{2!}, & \tau \leq t < 2\tau, \\
\vdots & \vdots \\
1 + \bar{k}_2 \frac{t}{\tau} + \bar{k}_2^2 \frac{(t-\tau)^2}{2!} + \cdots + \bar{k}_2^m \frac{[t-(m-1)\tau]^m}{m!}, & (m-1)\tau \leq t < m\tau, \\
\vdots & \vdots 
\end{cases}
\]
and \( e_{\tau}^{\bar{k}_2 t} \) is a solution to the following linear homogeneous equation with pure delay
\[
(3.22) \quad \begin{cases} 
\frac{dz}{dt} = \bar{k}_2 z(t-\tau), & t \geq 0, \\
z(s) \equiv 1, & s \in [-\tau, 0]. 
\end{cases}
\]

Lemma 3.8. Let \( k_1 \geq 0 \) and \( k_2 \geq 0 \). Then the solution \( z(t) \) to (3.20) (or equivalently (3.21)) satisfies
\[
|z(t)| \leq C_0 e^{-k_1 t} e_{\tau}^{\bar{k}_2 t},
\]
where
\[
C_0 := e^{-k_1 \tau} |z_0(-\tau)| + \int_{-\tau}^{0} e^{k_1 s} |z_0'(s) + k_1 z_0(s)| \, ds,
\]
and the fundamental solution \( e_{\tau}^{\bar{k}_2 t} \) with \( \bar{k}_2 > 0 \) to (3.22) satisfies
\[
e_{\tau}^{\bar{k}_2 t} \leq C (1 + t)^{-\gamma} e_{\tau}^{\bar{k}_2 t}
\]
for arbitrary number \( \gamma > 0 \).

Furthermore, when \( k_1 \geq k_2 \geq 0 \), there exists a constant \( \epsilon_1 = \epsilon(\tau) \) with \( 0 < \epsilon < 1 \) for \( \tau > 0 \), and \( \epsilon_1 = 1 \) for \( \tau = 0 \), and \( \epsilon_1 = \epsilon(\tau) \to 0 \) as \( \tau \to +\infty \), such that
\[
e^{-k_1 t} e_{\tau}^{\bar{k}_2 t} \leq C e^{-\epsilon_1 (k_1-k_2) t}, \quad t > 0,
\]
and the solution \( z(t) \) to (3.20) satisfies
\[
|z(t)| \leq C e^{-\epsilon_1 (k_1-k_2) t}, \quad t > 0.
\]
Taking Fourier’s transform to (3.9), and denoting the Fourier’s transform of $U^+(t, \xi)$ by $\hat{U}^+(t, \eta)$, we obtain

\begin{equation}
\tag{3.24}
\begin{aligned}
\frac{d}{dt}\hat{U}^+(t, \eta) + A(\eta)\hat{U}^+(t, \eta) &= B(\eta)\hat{U}^+(t - \tau, \eta), \\
\hat{U}^+(s, \eta) &= \hat{U}^+_0(s, \eta), \quad s \in [-\tau, 0], \quad \eta \in \mathbb{R},
\end{aligned}
\end{equation}

where

\begin{equation}
\tag{3.25}
\begin{aligned}
A(\eta) &= c_1 - d(e^{\lambda+i\eta} + e^{-(\lambda+i\eta)}) + ic\eta, \\
B(\eta) &= g'(0)\sum_{i \in \mathbb{Z}} e^{-i((i+c)\tau)\eta} K(i)e^{-\lambda((i+c)\tau)}.
\end{aligned}
\end{equation}

By the solution formula (3.21) in Lemma 3.7, the linear time-delayed ordinary differential equation (3.24) can be solved by

\begin{equation}
\tag{3.26}
\hat{U}^+(t, \eta) = e^{-A(\eta)(t+\tau)}e^{\overline{B}(\eta)t}\hat{U}^+_0(-\tau, \eta) \\
+ \int_{-\tau}^{0} e^{-A(\eta)(t-s)}e^{\overline{B}(\eta)(t-\tau-s)} \left[ \partial_s \hat{U}^+_0(s, \eta) + A(\eta)\hat{U}^+_0(s, \eta) \right] ds,
\end{equation}

where

\begin{equation}
\tag{3.27}
\overline{B}(\eta) := B(\eta)e^{A(\eta)\tau}.
\end{equation}

Taking the inverse Fourier transform to (3.26), we have

\begin{equation}
\tag{3.28}
U^+(t, \xi) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\xi\eta} e^{-A(\eta)(t+\tau)} e^{\overline{B}(\eta)t}\hat{U}^+_0(-\tau, \eta) d\eta \\
+ \frac{1}{2\pi} \int_{-\tau}^{0} \int_{-\infty}^{\infty} e^{i\xi\eta} e^{-A(\eta)(t-s)} e^{\overline{B}(\eta)(t-\tau-s)} \left[ \partial_s \hat{U}^+_0(s, \eta) + A(\eta)\hat{U}^+_0(s, \eta) \right] d\eta ds.
\end{equation}

We now estimate the decay rate of $U^+(t, \xi)$.

**Lemma 3.9.** Let the initial data $U^+_0(s, \xi)$ be such that

\[ U^+_0 \in C([-\tau, 0]; W^{1,1}(\mathbb{R})) \quad \text{and} \quad \partial_s U^+_0 \in L^1([-\tau, 0]; L^1(\mathbb{R})). \]

Then

\[ \|U^+(t)\|_{L^\infty(\mathbb{R})} \leq Ce^{-\mu_2 t} \]

for any $c > \tilde{c}$ and some $\mu_2 > 0$.

**Proof.** By Parseval’s inequality, it then follows from (3.26) that

\begin{equation}
\tag{3.29}
\|U^+(t)\|_{L^\infty(\mathbb{R})} \leq \|\hat{U}^+(t)\|_{L^1(\mathbb{R})} \leq I_1(t) + I_2(t),
\end{equation}

where $I_1(t)$ and $I_2(t)$ are given by

\begin{equation}
\tag{3.30}
I_1(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\xi\eta} e^{-A(\eta)(t+\tau)} e^{\overline{B}(\eta)t}\hat{U}^+_0(-\tau, \eta) d\eta,
\end{equation}

\begin{equation}
\tag{3.31}
I_2(t) = \frac{1}{2\pi} \int_{-\tau}^{0} \int_{-\infty}^{\infty} e^{i\xi\eta} e^{-A(\eta)(t-s)} e^{\overline{B}(\eta)(t-\tau-s)} \left[ \partial_s \hat{U}^+_0(s, \eta) + A(\eta)\hat{U}^+_0(s, \eta) \right] d\eta ds.
\end{equation}

By (3.24), we have

\begin{equation}
\tag{3.32}
\partial_s I_1(t) \leq -\mu_1 I_1(t),
\end{equation}

\begin{equation}
\tag{3.33}
\partial_s I_2(t) \leq -\mu_2 I_2(t),
\end{equation}

where $\mu_1$ and $\mu_2$ are positive constants. Integrating (3.32) and (3.33) from $0$ to $t$, we obtain

\begin{equation}
\tag{3.34}
I_1(t) \leq I_1(0) e^{-\mu_1 t},
\end{equation}

\begin{equation}
\tag{3.35}
I_2(t) \leq I_2(0) e^{-\mu_2 t},
\end{equation}

where $I_1(0)$ and $I_2(0)$ are bounded by the initial data.

Therefore, we have

\begin{equation}
\tag{3.36}
\|U^+(t)\|_{L^\infty(\mathbb{R})} \leq Ce^{-\mu_2 t},
\end{equation}

for some $c > \tilde{c}$ and $\mu_2 > 0$. 

where
\[ I_1(t) := \int_{-\infty}^{\infty} e^{-A(\eta)(t+\tau)} e_T(\eta) t \hat{U}_0^+(\tau, \eta) \, d\eta, \]
\[ I_2(t) := \int_{-\tau}^{0} \int_{-\infty}^{\infty} e^{-A(\eta)(t-s)} e_T(\eta) (t-\tau-s) \left[ \partial_s \hat{U}_0^+(s, \eta) + A(\eta) \hat{U}_0^+(s, \eta) \right] \, ds \, d\eta. \]

We are going to estimate \( I_i(t), i = 1, 2 \). By (3.25), one has
\[ |e^{-A(\eta)(t+\tau)}| = e^{-c_1(t+\tau)} \left| \exp \left( (t+\tau) d(e^{\lambda+\eta} + e^{-\lambda-i\eta}) \right) \right| \]
\[ = e^{-c_1(t+\tau)} \exp \left( (t+\tau) d(e^{\lambda} \cos \eta + e^{-\lambda} \cos \eta) \right) \]
\[ = e^{-c_1(t+\tau)} \exp \left( (t+\tau) d(e^{\lambda} + e^{-\lambda}) \cos \eta \right) \]
\[ = e^{-k_1(c)(t+\tau)} e^{-m(\eta)(t+\tau)}, \]

where
\[ k_1(c) := c_1 - d(e^{\lambda} + e^{-\lambda}) \quad \text{and} \quad m(\eta) := (1 - \cos \eta)d(e^{\lambda} + e^{-\lambda}). \]

By (3.27) and (3.25), we obtain
\[ |e_T(\eta)| = |B(\eta)||e^{A(\eta)\tau}| \leq e^{d(e^{\lambda}+e^{-\lambda})\tau} g'(0)G(\lambda)e^{k_1(c)\tau} =: \bar{k}_2(c), \]
where
\[ \bar{k}_2(c) := k_2(c)e^{k_1(c)\tau} \quad \text{and} \quad k_2(c) := e^{d(e^{\lambda}+e^{-\lambda})\tau} g'(0)G(\lambda). \]

Then we get
\[ |e_T(\eta)|^t \leq e^{\bar{k}_2(c)t}. \]

It is easy to see that
\[ k_1(c) - k_2(c) = c \lambda - d(e^{\lambda} + e^{-\lambda} - 2) + 1 - e^{d(e^{\lambda}+e^{-\lambda})\tau} g'(0)G(\lambda) \]
\[ = c \lambda - d(e^{\lambda} + e^{-\lambda} - 2) + 1 - e^{d(e^{\lambda}+e^{-\lambda})\tau} g'(0) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda(i+c\tau)} \]
\[ \geq c \lambda - d(e^{\lambda} + e^{-\lambda} - 2) + 1 - e^{d(e^{\lambda}+e^{-\lambda})\tau} g'(0) \sum_{i \in \mathbb{Z}} K(i)e^{-\lambda i} > 0, \]
due to (2.5). Let \( \mu_0 := k_1(c) - k_2(c) > 0 \). Then by (3.23), we have
\[ |e^{-A(\eta)(t+\tau)} e_T(\eta)| \leq e^{-m(\eta)(t+\tau)} e^{-k_1(c)(t+\tau)} e^{k_2(c)t} \]
\[ \leq C e^{-\varepsilon_1[k_1(c)-k_2(c)]} e^{-m(\eta)t} = C e^{-\varepsilon_1\mu_0 t} e^{-m(\eta)t}. \]

Applying (3.29), we derive the optimal estimate for \( I_1(t) \):
\[ I_1(t) = \int_{-\infty}^{\infty} e^{-A(\eta)(t+\tau)} e_T(\eta) t \hat{U}_0^+(\tau, \eta) \, d\eta \]
\[ \leq C e^{-\varepsilon_1\mu_0 t} \| \hat{U}_0^+ \|_{L^\infty(\mathbb{R})} \int_{-\infty}^{\infty} e^{-m(\eta)t} \, d\eta \]
\[ \leq C e^{-\mu_2 t} \| U_0^+ \|_{L^1(\mathbb{R})}, \]
where \( \mu_2 := \varepsilon_1 \mu_0 \).

Similarly, we can estimate \( I_2(t) \). Note that
\[
\sup_{\eta \in \mathbb{R}} |A(\eta)\hat{U}_0^+(s, \eta)| = \sup_{\eta \in \mathbb{R}} \left| [c_1 - d(e^{\lambda+4i\eta} + e^{-(\lambda+4i\eta)}) + ic\eta] \hat{U}_0^+(s, \eta) \right| \leq C\|U_0^+(s)\|_{W^{1,1}(\mathbb{R})}.
\]

Thus, we can derive the decay rate for \( I_2(t) \) as follows:
\[
I_2(t) = \int_{-\tau}^{0} \int_{-\infty}^{\infty} \left| e^{-A(\eta)(t-s)} e^{B(\eta)(t-\tau-s)} \left[ \partial_x \hat{U}_0^+(s, \eta) + A(\eta) \hat{U}_0^+(s, \eta) \right] \right| \, d\eta ds \\
\leq C \int_{-\tau}^{0} \int_{-\infty}^{\infty} e^{-\varepsilon_1 \mu_0(t-s)} e^{-m(\eta)(t-s)} \left| \partial_x \hat{U}_0^+(s, \eta) + A(\eta) \hat{U}_0^+(s, \eta) \right| \, d\eta ds \\
\leq Ce^{-\varepsilon_1 \mu_0 t} \int_{-\tau}^{0} \int_{-\infty}^{\infty} e^{-m(\eta)(t-s)} \sup \left| \partial_x \hat{U}_0^+(s, \eta) + A(\eta) \hat{U}_0^+(s, \eta) \right| \, d\eta ds \\
\leq Ce^{-\varepsilon_1 \mu_0 t} \left[ \|\partial_x U_0^+(s)\|_{L^1([-\tau,0];L^1(\mathbb{R}))} + \|U_0^+(s)\|_{L^1([-\tau,0];W^{1,1}(\mathbb{R}))} \right].
\]

Substituting (3.30) and (3.31) to (3.28), we obtain
\[
\|U^+(t)\|_{L^\infty(\mathbb{R})} \leq Ce^{-\mu_2 t} \quad \text{for } c > \tilde{c}.
\]

The proof is complete. \( \square \)

Let us choose \( U_0^+(s, \xi) \) such that \( U_0^+ \in C([-\tau,0];W^{1,1}(\mathbb{R})) \) and \( \partial_x U_0^+ \in L^1([-\tau,0];L^1(\mathbb{R})) \), and \( U_0^+(s, \xi) \geq |\hat{U}_0(s, \xi)| \) for \( (s, \xi) \in [-\tau,0] \times \mathbb{R} \). Then combining Lemmas 3.6 and 3.9, we immediately obtain the convergence rate for \( \tilde{U}(t, \xi) \).

**Lemma 3.10.** When \( \tilde{U}_0 \in C([-\tau,0];W^{1,1}(\mathbb{R})) \) and \( \partial_x \tilde{U}_0 \in L^1([-\tau,0];L^1(\mathbb{R})) \), then
\[
\|\tilde{U}(t)\|_{L^\infty(\mathbb{R})} \leq Ce^{-\mu_2 t} \quad \text{for } c > \tilde{c}.
\]

**Proposition 3.11.** It holds
\[
\sup_{\xi \in (-\infty, x_0]} |U(t, \xi)| \leq Ce^{-\mu_2 t} \quad \text{for } c > \tilde{c}.
\]

**Proof.** Notice that \( \tilde{U}(t, \xi) = \sqrt{\omega(\xi)}U(t, \xi + x_0) = e^{-\lambda \xi}U(t, \xi + x_0) \), and \( \sqrt{\omega(\xi)} = e^{-\lambda \xi} \geq 1 \) for \( \xi \in (-\infty, 0] \). Then we obtain
\[
\sup_{\xi \in (-\infty, 0]} |U(t, \xi + x_0)| \leq \|\tilde{U}(t)\|_{L^\infty(\mathbb{R})} \leq Ce^{-\mu_2 t} \quad \text{for } c > \tilde{c},
\]

which implies
\[
\sup_{\xi \in (-\infty, x_0]} |U(t, \xi)| \leq Ce^{-\mu_2 t} \quad \text{for } c > \tilde{c}.
\]

The proof is complete. \( \square \)
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