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ABSTRACT

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GASTOF: Ultra-fast ToF forward detector for exclusive processes at the LHC *

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GASTOF (Gas Time-of-Flight) detector is a Cherenkov detector proposed for very precise ($\delta t \sim 10–20 \text{ ps}$) arrival time measurements of forward protons at some 420 m from the central detectors of CMS and ATLAS. Such an excellent time resolution will allow by z-by-timing technique for precise measurement of the z-coordinate of the event vertex in exclusive production at the LHC, when two colliding protons are scattered at very small angles. In the paper we present first GASTOF prototype, simulations of its performance as well as first tests using a cosmic muon telescope.

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1. Introduction

In the central exclusive processes $pp \rightarrow p \oplus X \oplus p$, where $\oplus$ denotes the absence of hadronic activity, the central system $X$ can be produced either via $\gamma \gamma$ fusion or in exclusive diffractive production. Detection of the forward scattered protons will allow for new and complementary studies at the LHC. Already at low luminosity, the two photon dilepton production, precisely known from QED, can serve as a calibration process [1]. For few fb$^{-1}$ of the integrated $pp$ luminosity the high energy photon physics opens up giving access to precision studies of quartic gauge couplings, anomalous $W$ or $Z$ pair production, and at higher luminosities, super-symmetric particle pair production in very clean environment [2]. Starting from tens of fb$^{-1}$ the exclusive diffractive production of the Higgs boson becomes important [3, 4].

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There are three main reasons why the exclusive diffractive production is important especially for the Higgs boson studies at the LHC:

1. Forward scattered protons, tend to "select" state of the central system $X$ to be $J_z = 0$, $C$ and $P$ even ($0^{++}$). Moreover, correlations between outgoing proton momenta are sensitive to these quantum numbers.

2. Mass resolution of about 2% of the central system can be achieved from momenta measurements of the scattered protons alone, which would help to resolve a nearly degenerate super symmetric Higgs sector, for example [5].

3. Excellent signal to background ratio of about unity for the SM Higgs production and more than an order of magnitude larger for certain MSSM scenarios.

The FP420 collaboration proposes [6] to install proton detectors at some 420 m from the interaction point (IP) of CMS or ATLAS. Acceptance of such detectors matches very well energy distributions of the forward protons in the exclusive production of the light Higgs boson. As the cross section for diffractive production of the SM Higgs production is expected to be small, 1–10 fb depending of the Higgs boson mass [4], it is imperative to measure it at high LHC luminosity with many interactions per beam crossing (event pile-up).

The aim of a very precise time measurement using GASTOF [1] is to determine $z$-coordinate of event vertex using the $z$-by-timing technique, and consequently to match it with the vertex measured by central detectors. The $z$-by-timing technique is based on the arrival time difference for two protons detected on both sides of the IP. If detectors of forward protons are at distance $L$, and the event vertex is displaced from the nominal IP at $z = 0$ to some $z$ (width of the longitudinal distribution of the interaction point is expected to be of 50–70 mm) then

$$\Delta t = \frac{L + z}{c} - \frac{L - z}{c} = \frac{2z}{c}$$

(1)

where $\Delta t$ is the arrival time difference of the two protons to be measured with GASTOF detectors, and $c$ is speed of light ($\beta \approx 1$ for the forward protons). It follows from Eq. (1) that precision of $z$ measurement is $\delta z = c\delta t/\sqrt{2}$. Hence the detector time resolutions of $\delta t = 20$ or 10 ps would result in 4 mm or 2 mm resolutions in $z$, respectively. For studies of exclusive processes at high LHC luminosity this will be essential to reduce accidental coincidences due to event pile-up, where the two forward protons and the central system $X$ are not coming from the same interaction.

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1 Or, using a complementary QUARTIC detector which is based on quartz radiator.
2. GASTOF

GASTOF is a Cherenkov gas detector – photons produced by high energy protons traversing gas medium are reflected by a mirror onto a very fast photomultiplier. In gases, thanks to small refractive index, the Cherenkov photons are radiated at very small angles, and propagate at speed very close to c, therefore very good time resolution is expected. The detector is simple, very robust, and light – the multiple scattering induced by GASTOF should be small and allow for placing it in front of or between the planes of the proton tracking detectors, without affecting the eventual resolutions. The FP420 forward detector sensitive areas will be heavily irradiated, in particular by single diffractive protons produced at huge rates at the LHC – if needed the gas can be flushed therefore it is expected that GASTOF will be radiation hard. In addition, thanks to its directionality and relatively high energy threshold of incident particles GASTOF will not be too sensitive to stray charged particles in the LHC tunnel. On the other hand, in gas the Cherenkov photon yield is not large, and the radiator length is limited – therefore, as the photon spectrum is peaked at short wavelengths, a lot of effort is put into providing good efficiency of detecting UV photons.

First two identical GASTOF prototypes have been built to test these expectations. A 31 cm long, 6 cm square tube, is filled with C\textsubscript{4}F\textsubscript{10} at 1.2 atm, with refractive index n\textasciitilde1.0018. A flat mirror at 45\degree reflects Cherenkov light onto a 2 inch square photocathode of the micro-channel plate photomultiplier (MCP-PMT) 85011-501 from Burle \cite{7}). Special UV coated mirrors have been used, which have non-zero reflectivity for \( \lambda > 160 \) nm and more than 75\% above 180 nm. The Burle MCP-PMTs have UV grade fused silica windows, the collection efficiency of 50\%, and multiple anodes in form of 8x8 anode matrix. They are characterized by a sharp rise time of 300 ps and low transit time jitter of about 40 ps.

A simple, based on ray-tracing, Monte Carlo simulation has been prepared, and its results for the prototype are shown in the Fig.\cite{1}. For high energy charged particles hitting GASTOF centrally, along its axis, almost 200 photons in average are radiated and hit the mirror. This results in about 13 photoelectrons produced in average at the photocathode. The light spot at the photocathode has diameter of 4 cm. Finally, all the Cherenkov photons arrive at the photocathode within a 4 ps time window! These results indicate that indeed GASTOF can provide efficient, and extremely fast and accurate timing signal at the LHC.

3. First Results

The GASTOF test stand has been prepared using a simple cosmic ray telescope, as sketched in Fig.\cite{2}. Two small plastic scintillator blocs sepa-
Fig. 1. Results of simulation of the GASTOF prototype – 31 cm long, filled with C₄F₁₀ at 1.2 atm, with refractive index \( n \approx 1.0018 \). One expects that in average, for one high energy proton hitting GASTOF centrally, about 198 Cerenkov photons hit the mirror, and 13.4 photoelectrons are produced at the Burle MCP-PMT. Upper, left plot shows spatial distribution of the photons at the photocathode. Upper, right plot shows arrival time of these photons, where \( t = 0 \) is set for a proton entering GASTOF. Lower plots show, respectively, numbers of produced photons and photoelectrons as a function of the photon wavelength. Each plot shows number of events per bin, hence sums of the bin contents are equal to the total number of photons and photoelectrons, respectively.

rated vertically and readout by Philips PMTs XP2020 are used as a cosmic muon trigger. In each MCP-PMT the 4x4 central group of anodes was connected together by short wires of equal lengths, and the rest of anodes was grounded. The signals from the Burle MCP-PMTs are sent via about 20 cm long SMA cables to very fast Hamamatsu C5594 amplifiers. Two GASTOF prototypes placed one after the other and were tested simultaneously. Two GASTOF signals as well as those from the trigger are read using a fast
3GHz LeCroy Wavepro 7300A scope with digital resolution of 50 ps \[8\].

Fig. 2. Cosmic ray test stand for the GASTOF prototypes. Plastic scintillators (up and down) readout by Philips photomultiplier tubes XP2020 are used as a trigger. The signals from the Burle MCP-PMTs are sent via short SMA cables to very fast Hamamatsu C5594 amplifiers. Signals from the GASTOF prototypes as well as from the trigger are read using a fast 3GHz LeCroy Wavepro 7300A scope.

In Fig. (3) an example of cosmic ray signals from two GASTOF detectors is shown. More than hundred such events were collected in a one day run, allowing to make first statistical analysis of the data. The signals were fitted using the Landau distribution function, and the crucial parameters were extracted – the arrival and rise time of each pulse. The average rise time is similar for both detectors and is about 600 ps. This is two times worse than the single anode rise time quoted by Burle, therefore it is believed to be worsened by the anode grouping. The time difference distribution was also measured – it is of gaussian shape with about 100 ps width. Assuming that this width is dominated by the time resolution of two prototypes, and that they are the same, the upper limit of 70 ps on a single detector resolution can be set. This is about two times bigger than the transit time jitter expected for the Burle MCP-PMT, but is consistent with the degradation of the Burle rise time, possibly due to the anode grouping.

Results from these first tests are encouraging, showing a big potential of GASTOF detectors in domain of ultra-precise timing applications. Next
Fig. 3. An example of cosmic ray signals from two GASTOF detectors. The data were taken with the LeCroy Wavepro 7300A scope with digital resolution of 50 ps. The Landau distribution function was used for fits.

steps will include use of the single anode readout or improved anode grouping scheme, as well as of new Burle MCP-PMTs with even smaller transit time jitter.

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