Modelling and Assessment of Accidental Gas Release from Damaged Subsea Pipelines

Ahmed M. Ellethy, Ahmed S. Shehata, Ali I. Shehata, and Ahmed Mehanna

Abstract—Due to the high demand on hydrocarbon exploration nowadays and in future overall the world, the risk potential of Subsea gas and oil leakage increases and can lead to a catastrophic incidents such as explosion, fire or loss stability of subsea constructions. Therefore, the purpose of this paper is to assess and control the potential of hazards for the gas flows from subsea pipelines rupture to sea surface with varying of influence parameters on gas plume. 2D Computational Fluid Dynamics (CFD) simulations have been carried out. The influence parameters that contained in the present paper leakage hole sizes and water depths. The Rotvoll experiment data is used in order to validate numerical computational fluid dynamics model. The subsea gas diffusion behavior is investigated for the risk assessment aspect. The main parameters which included to assess the gas dispersion are rising time and fountain height. Observed from our sceneries results that the water depth has a major effect on gas rise time where the deeper water depth is about 3 times longer than the lower water depth to reach water surface whilst the hole size has a high effect on the fountain height where the larges hole size is approximately 2.1 times higher than the smaller hole size. So our objective is helping the petroleum companies to find a mitigation method by modelling and assessing the gas release in order to contain the subsea gas leakage.

Index Terms—Underwater gas leakage, risk assessment, gas dispersion, hole size, water depth.

I. INTRODUCTION

After centuries of technology the exploration and extraction of hydrocarbon resources have expanded offshore and into deeper waters, the Gulf of Mexico is one of the most important regions for energy resources and infrastructure for both onshore and offshore. Gulf of Mexico's offshore oil production reached to 17% of total U.S. oil production. Offshore gas production in the Gulf contributed for 5% of U.S. dry production in total. Over than 45% of U.S. oil refining capacity was placed along the Gulf Coast [1]. As seen in Fig. 1. Therefore, the risk of potential leakage is also rising at the same time with the number of subsea structure installations and pipelines increases. Subsea extraction and transport techniques have been developed, including the subsea pipeline. Pipelines are widely used to transport the hydrocarbon fluids and fossil fuels and chemicals over millions of miles across the world.

The pipeline structures are designed to Steadfastness against several severe environmental conditions in order to maintain an environment and reliable distribution along the pipeline production from the production wells to shore. And pipeline considers one of the most important assets for petroleum companies. Likewise, the more professional approach by developing the Risk assessment plans for any offshore structures installations, including the predicting impact consequence of a variety of catastrophic events. This could involve the leakage from any hydrocarbon containment from topsides or subsea equipment such as subsea pipelines, risers, production x-tree and production manifold [2].

The sub-sea pipeline release relies on if the leakage fluids are in liquid or gaseous form. If the leakage is in the liquid phase, the buoyancy will contribute in the diffusion of the liquid due to the liquid release is lighter than the water on the surface of the water. This diffusion will lead to pollution spot, or it will lead to a pool of fire. In the case of the ignition source is exist, Even though the gas buoyancy is more larger than for liquids leakage, drag forces will allow the plume to split up between the leakage point and the water surface, and the released gas will rise to the surface as a stream of gas bubbles [3].

Many researchers have implemented a number of studies on the leakage of underwater gas in the last decades and have made a significant contribution. Many factors have been shown to have impacts on underwater plume behavior, such as sea currents, gas dissolution, hydrate creation and hydrate decomposition by [4].

Li. et al. [5] Studied the systematic simulation and evaluation of subsea release gas deflagration and diffusion using an Integrated CFD model to analyze the development of the incident scenario and The jack-up drilling rig is used as a case in point to identify the effects of gas deflagration and diffusion on surface platform. Xu. et al. [6] investigated a systematic simulation and assessment of gas dispersion above sea from a Subsea gas release by developing a 3D CFD model to evaluate the behavior patterns of flammable gas above sea level, and the jack-up drilling rig is used to demonstrate the effect of the cloud of gas on surface. The

Fig. 1. The Gulf of Mexico area modified from [1].
author has observed that the size of the hazardous area decreases with an increase in the release rate of the surface gas. Also Li. et al. [7] presented an experimental and numerical analysis of the release and dispersion behavior of subsea gas. A variety of release scenarios are conducted to study the impact of release pressure, water depth and leak direction on dispersion behavior. The author observed that the gas plume under the larger leak pressure has been shown to have a higher volume. The effect of the leaking pressure on the gas plume is also equivalent to the size of the nozzle. An experimental has been designed in order to test gas leakage rate and dispersion behavior in various scenarios which is getting from (3).

The mass conservation equation is shown as in (1), where \( \alpha \) is volume fraction of \( q_{th} \) phase, \( \rho_f \) is the density of \( q_{th} \) phase and \( \bar{v}_q \) is the velocity of \( q_{th} \) phase. 

\[
\frac{\partial}{\partial t} (\alpha \rho_f) + \nabla \cdot (\alpha \rho_f \bar{v}_q) = 0 
\] (1)

The momentum equation is shown as in (2), where \( \bar{v}_q \) is the velocity of \( q_{th} \) phase, \( \rho \) is the density of mixed phases, which is getting from (3), \( \rho_f \) and \( \bar{f} \) are the gravitational body force and external body force, \( \mu \) is viscosity that is the sum of molecular mixture viscosity and turbulent viscosity. So the characteristics \( \rho \) and \( \mu \) are the most important as the momentum equation rely on volume fractions of phases.

\[
\frac{\partial}{\partial t} (\rho \bar{v}) + \nabla \cdot (\rho \bar{v} \bar{v}) = -\nabla p + \nabla \cdot (\mu \nabla \bar{v} + \nabla \bar{v}^T) + \rho \bar{g} + \bar{f} 
\] (2)

### II. METHODOLOGY

**A. Modeling Approach of Underwater Gas Release**

The Eulerian-Lagrangian modelling approach is a hybrid two-way coupling modeling approach consisting of the Eulerian-Lagrangian Discrete Phase Model (DPM) and the Eulerian-Eulerian Volume of Fluid Model (VOF) also the common name is coupled DPM and VOF model. In this paper the discrete particles are referred to the bubble and the continuous phase is referred to the water (The fluid in the computational domain). DPM tracks the movement of discrete particles, whilst VOF describes the movement of the continuous phase. The volume ratio is measuring parameter for the distribution of different phases. In the computational domain, the Eulerian volume of fluid model considers the volume occupied by the released gas and the summing volume of the water and the released gas equal one.

**B. Theoretical Model**

1) The Eulerian-Eulerian volume of Fluid Model (VOF)

VOF is a strategy of free surface tracking under the constant Eulerian mesh. The interacting mediums are all different phases in the computational domain and the volume of one phase cannot be occupied by other phases. The governing equations show the interakt between discrete particles move through continuous phase by exchanging mass conservation, momentum conservation, and energy conservation [12].

The mass conservation equation is shown as in (1), where \( \alpha_q \) is volume fraction of \( q_{th} \) phase, \( \rho_f \) is the density of \( q_{th} \) phase and \( \bar{v}_q \) is the velocity of \( q_{th} \) phase.

\[
\frac{\partial}{\partial t} (\alpha_q \rho_f) + \nabla \cdot (\alpha_q \rho_f \bar{v}_q) = 0 
\] (1)

The momentum equation is shown as in (2), where \( \bar{v}_q \) is the velocity of \( q_{th} \) phase, \( \rho \) is the density of mixed phases, which is getting from (3), \( \rho_f \) and \( \bar{f} \) are the gravitational body force and external body force, \( \mu \) is viscosity that is the sum of molecular mixture viscosity and turbulent viscosity. So the characteristics \( \rho \) and \( \mu \) are the most important as the momentum equation rely on volume fractions of phases.

\[
\frac{\partial}{\partial t} (\rho \bar{v}) + \nabla \cdot (\rho \bar{v} \bar{v}) = -\nabla p + \nabla \cdot (\mu \nabla \bar{v} + \nabla \bar{v}^T) + \rho \bar{g} + \bar{f} 
\] (2)

| Authors | Numerical models | Turbulence model | Hole size | Water depth | Validation model | Pressure discretization | Continuity, momentum and turbulence equations | Pressure velocity coupling | Volume fraction |
|---------|------------------|------------------|-----------|-------------|------------------|------------------------|---------------------------------|----------------------------|----------------|
| [9]     | Coupled VOF and DPM approach | The standard K-\( \varepsilon \) model | - | 30 m, 100 m and 400 m | Rotvoll experiment | PRESTO! | Second order upwind scheme | PISO scheme | The Geo-Recon struct algorithm |
| [10]    | Coupled VOF and DPM approach | The realizable K-\( \varepsilon \) turbulence model | 10 mm, 50 mm, 100 mm and full rupture | Several depth approximatively from 50 m to 235 m | Rotvoll experiment | PRESTO! | Second order upwind scheme | PISO scheme | The Geo-Recon struct algorithm |
| [11]    | Coupled VOF and DPM approach | The realizable K-\( \varepsilon \) turbulence model | 30 mm | 30 m, 50 m, and 70 m | Rotvoll experiment | PRESTO! | Second order upwind scheme | PISO scheme | The Geo-Recon struct algorithm |
| [12]    | Coupled VOF and DPM approach | The standard K-\( \varepsilon \) model | 25.4 mm | 50 m and 300 m | Rotvoll experiment | PRESTO! | Second order upwind scheme | PISO scheme | The Geo-Recon struct algorithm |
| [13]    | Coupled VOF and DPM approach | The standard K-\( \varepsilon \) model | - | 7 m | Rotvoll experiment | PRESTO! | Second order upwind scheme | PISO scheme | The Geo-Recon struct algorithm |

**TABLE I: INFORMATION ABOUT THE PREVIOUS STUDIES**

1) Li., et al. [7] presented an experimental and numerical analysis of the release and dispersion behavior of subsea gas. 2) Rotvoll et al. [9] presented an experimental and numerical analysis of the release and dispersion behavior of subsea gas. 3) Li., et al. [8] observed that the size of the hazardous area decreases with an increase in the release rate of the surface gas. 4) Li. et al. [7] observed that the gas release rate increased approximately linearly with an increasing in pressure at 0.75 m. 5) Li., et al. [8] observed that the gas release rate increased approximately linearly with an increasing in pressure at 0.75 m.

Therefore, the purpose of this study is to investigate the behavior of gas release and dispersion characteristics and also to determine subsea rise time, gas transport path and fountain height until reach sea surface with using production gas density of Petroleum Egyptian company around 0.648 kg /m³ and study the effects of the hole size and water depth on subsea gas release by developing a CFD model in order to find mitigation method to contain the subsea gas release.
\[ P = \sum a_{q} \rho_{q} \]  

(3)

The energy equation is shown as in (4), where \( k_{\text{eff}} \) is the effective thermal conductivity coefficient; \( Sh \) is the source term which is equal 0 by default.

\[ \frac{\partial}{\partial t} (\rho E) + \nabla \cdot (\rho \mathbf{v}) = \nabla \cdot (k_{\text{eff}} \nabla T) + Sh \]  

(4)

The standard k-epsilon model is the model of turbulence. In order to represent the turbulent properties of the flow, the standard k-epsilon model is comprised of two extra transport equations. Turbulent kinetic energy, the first transport parameter, \( K \), the second transport parameter is turbulent dissipation, \( \varepsilon \), the two equations are shown below in (5) & (6).

\[ \frac{\partial (\rho K)}{\partial t} + \nabla \cdot (\rho \mathbf{v} K) = \nabla \cdot (\mu_{t} \nabla K) + 2\mu_{t} S_{ij} S_{ij} - \rho \varepsilon \]  

(5)

\[ \frac{\partial (\rho \varepsilon)}{\partial t} + \nabla \cdot (\rho \mathbf{v} \varepsilon) = \nabla \cdot (\mu_{t} \nabla \varepsilon) + C_{1\varepsilon} \frac{\varepsilon}{K} 2\mu_{t} S_{ij} S_{ij} - C_{2\varepsilon} \rho \varepsilon \]  

(6)

where \( \sigma_{k}=1.00 \), \( \sigma_{\varepsilon}=1.30 \), \( C_{2\varepsilon}=1.92 \), \( C_{1\varepsilon}=1.44 \), and The flow was believed to be fully turbulent, and the molecular viscosity influences were negligible. Therefore, the standard k-\( \varepsilon \) model is valid for completely turbulent flows only. This model’s drawback is delayed and separation decreased.

2) Eulerian-Lagrangian discrete phase model (DPM)

DPM tracks the movement of discrete particles (bubbles) in the water from the released gas. The trajectory of the bubbles is solved by the particle force equation in the differential Lagrangian coordinate frame. In Cartesian coordinates frame, the force balance equation of particles is shown in (7). Bubble inertia is proportional to the different forces that act with bubbles, additional to drag force, buoyancy and mass force [14].

\[ \frac{du_{b}}{dt} = \frac{g(\rho_{b} - \rho)}{\rho_{b}} + F_{D} + F_{VM} \]  

(7)

The subsea gas plume motion is affected by the changing in Temperature and Hydrostatic Pressure Losses. Temperature and hydrostatic pressure are the main factors for bubble properties such as size and density, the ideal gas law usually defines them. The movement of bubbles is combined with the continuous phase flow. To measure the flow of the continuous phase, an Eulerian VOF model preserving momentum and mass via the Navier-Stokes equations. The continuous phase mass transformation is represented by the continuity equation, which is a simplified form considering that the flow is diluted and the impact of the bubbles can be ignored [15].

where \( u_{b} \) is the velocity of the bubbles, \( \rho_{b} \) is the density of the bubbles gas, \( u \) is the velocity of the fluid, \( \rho \) is the density of continuous phases, \( F_{D} \) is drag force, \( N \); \( F_{VM} \) is virtual mass force and determine the drag force from (8).

\[ F_{D} = \frac{18}{\rho_{b} a^{2}} \frac{C_{D} R e}{24} (u_{b} - u) \]  

(8)

where \( C_{D} \) is the drag coefficient, \( \rho_{b} \) is the density of the bubble gas, \( Re \) is the Reynolds number and \( (u_{b} - u) \) is a velocity difference between bubble velocity and liquid velocity, also the turbulent dispersion is generated by the turbulent velocity variation of the drag force. To evaluate the turbulent dispersion of gas bubbles, the Eulerian-Lagrangian model utilize a random tracking model so the turbulent velocity variance is calculated by (9).

\[ u' = \xi \sqrt{\frac{K}{\varepsilon}} \]  

(9)

where \( \xi \) is a random number subject to Gaussian distribution; \( K \) is turbulent kinetic energy.

In the K-\( \varepsilon \) model, the determined time of bubble particles moving in turbulent vortices as in (10).

\[ t_{\tau} = 0.15 \frac{K}{\varepsilon} \]  

(10)

where \( t_{\tau} \) is an existence period of turbulent vortices and \( \varepsilon \) turbulent energy dissipation rate.

C. Model Descriptions

1) Geometry and mesh model

The two dimensional geometry built in ANSYS design modeler the same dimensions as the experimental basin, with width of 9 m and depth of 7 m plus 3 m for air above the water surface. A uniform grid size is created in ANSYS Meshing then the mesh has been refined by region adapting in ANSYS Fluent, around the gas plume area is refined as shown in Fig. 2. (a) and (b). Increased number of cells till reached to the optimum grid cells that didn’t make more changes in simulation results, therefore, the simulations have performed with 96899 quadrilateral cells also used time step size at 0.001 to get better results and keep courtant number below 1.

2) Solution method

An unsteady solver based on pressure is used. For unsteady simulations, the PISO algorithm is sufficient and is used for pressure and velocity coupling [14]. In order to enhance model stability in the gravity field, the implicit body force formulation is enabled. For interface tracking of various phases, the Geometric-Reconstruct scheme is used. Stochastic monitoring (random walk) considers the additional drag force of the dispersion phase (gas bubbles) in the turbulent dispersion [9]. First order implicit for transient formulation and second order upwind for turbulent kinetic energy, momentum and turbulent dissipation rate. The numerical model is implemented by the fluent software with a set of UDFs. The viscosity and density of the water phase are set to 1.003e-3 kg/m.s and 998.23 kg/m³ respectively. For air, it was given as 1.7894e-5 kg/m.s and 1.225 kg/m³.

3) Model validation

A numerical model is generated with the same size as The experiment is called “Rotvoll” and summarized in the report.
of [16] to check the validity of the numerical model. A series of experiments were carried out in a rectangular basin with a depth of 7 m and 9 m width. The pipeline was linked at the release point between this vessel and the bottom centre of the basin. To hold the downstream pressure constant at a preset value to achieve a constant flow rate during release, a pressure regulator was installed in the pipeline. The underwater gas plume characteristics were captured by a high-speed video camera. The basin was filled with water and pure air released from the bottom with flow rate of 83, 170 and 750 Nl/s. The inlet consisted of a release valve with a fast acting piston that injects gas vertically with a 10 streams of particles are arranged at 0.33 m in front of the release point.

III. RESULTS AND DISCUSSION

A. Validation Results

In order to compare experimental and simulation findings, the rise time, centerline velocity and fountain height are chosen as reference parameters. There is some lack of experimental data for flow rate 0.75 Nm³/s. The simulation and experimental rise time are concise in Fig. 3. The rise time defines with the first particle reach the water surface. The simulation and experimental fountain height are summarized in table II and finally the Centre line velocity is summarized in Fig. 4. (a) and (b). For comparing experimental and simulation outcomes. The rise time and fountain height are selected as reference variables the same as [11]. Found that the outcomes of the simulation are magnificently aligned with the experiments, It can be shown that the simulated values are greater than the experimental values induced by the k– model, which only provides the average turbulence knowledge and lacks the ability to dissipate the elevated momentum around the gas core to water. The URANS LES turbulence model has been shown to display a more reliable estimate of plume centerline velocity than the RANS k– model. Furthermore, a disadvantage of overlong measurement time, which cannot be neglected. The current numerical model is still used for further discussion because the prediction error with the current turbulence model is inside a reasonable range and also the affordable time cost. The contour plot of volume fractions in Fig. 5 shows the deflected rising water upward in radial surface flow occurs due to the momentum of the entrained water plume, this elevation is clear evidence of the two-way coupling used in the present CFD simulation. The blue color represents the phase of the water (continuous phase) and the ambient air is represented by the red color. The yellow line reflects the interface between the two phases that illustrated by the Geo Reconstruction Scheme. Fig. 6 shows the bubble contour for particle residence time after 5 seconds at flow rate 0.17 Nm³/s.

| Flow rate (Nm³/s) | Fountain height (m) |
|------------------|---------------------|
| Experiment       | 0.083               |
| Simulation       | 0.25                |
|                  | 0.17                |
|                  | 0.56                |
|                  | 1.21                |
|                  | 0.75                |

Fig. 3. Rise time comparison for gas volume flow rates of 0.085 Nm3/s, 0.17 Nm³/s and 0.75 Nm³/s between experiment and simulation.

(a) Gas volume flow rate 0.085 Nm3/s

(b) Gas volume flow rate 0.17 Nm³/s

Fig. 4. Center line velocity comparison between experiment and simulation for gas volume flow rates (a) 0.085 Nm³/s and (b) 0.17 Nm³/s.

Fig. 5. Contour plot colored by the volume fractions at the water surface.

Fig. 6. Shows the bubble contour for particle residence time after 5 seconds at flow rate 0.17 Nm³/s.
B. Effect of Leakage Hole Size

The following scenarios with four subsea gas release rates corresponding to variety of hole sizes of 25 mm, 40 mm, 60 mm, 90 mm are chosen to evaluate the effect of the gas release rate of the underwater gas dispersion behavior on the same basin of validation model. The initial momentum and buoyancy are the main factors that control the growth of the gas plume, related to the lower gas release rate example, the initial release momentum is lower, therefore, leads to a longer increasing at the rising gas plume time. The comparison of gas plumes factors under various release rates is provided, suggesting that the rate of gas release has a major impact on the gas dispersion. Specifically, at the gas rising time were found that the increasing in a gas release rate leads to decrease in gas rise time. The smaller hole size 25 mm lead to more rise time 6.093 s and fountain height 0.32 m compared with hole size 90 mm the rise time 4.72 s and maximum fountain height 0.7 m. The growing trend of vertical rise distance with gas releasing time under various of hole sizes is shown in Fig. 7 and Fig. 9 shows the volume fraction contour for gas flow with water surface. It is apparent that the rate of gas leakage has significantly impacted on the growing gas plume. Therefore, the increase in the gas leakage rate, the lower rise time is expected. As can be concluded comparison between various scenarios of hole sizes with rise time of gas leakage and maximum fountain height in Fig. 8.

C. Effect of Water Depth

The depth of water is one of the main factors that impact on the rise time of gas plumes, so it’s important to clarify the effect of water depth and it should have a major role in risk assessment. The simulations have conducted by using four different depths as in shallow water (4 m, 7 m, 10 m and 13 m) with gas leakage rate (2 m/s). Fig. 10 presents the difference comparison among many parameters that got affected by water depth. Observed that the higher water depth 13 m lead to more rise time 9.91 s and fountain height 0.45 m compared with the smaller water depth 4 m the rise time 3.24 s and maximum fountain height 0.55 m. Fig. 11 shows the vertical rise distance with gas releasing time under a variety of water depths, its seem close curves due to the near velocity values in shallow depths. It’s clear to say the rise time increases and fountain height decreases with increasing the water depth. Volume fraction contour to illustrate the gas flow with water surface at variety of water depths are concluded in Fig. 12.
well notice that the higher water depth contributes to increase the gas plume diameter with considering it is not our topic in this paper.

![Graph showing rise time and fountain height vs. water depth](image)

Fig. 10. Shows the comparison among many parameters that got affected by water depth (a) Rise time; (b) Fountain height.

![Graph showing vertical rise distance with gas releasing time under a variety of water depths](image)

Fig. 11 The vertical rise distance with gas releasing time under a variety of water depths.

### IV. CONCLUSION

A numerical analysis was conducted based on CFD model to study the properties of underwater gas release and to assess the gas dispersion by focus on rising time and fountain height of gas plume in order to mitigate solution for petroleum companies. The main results are summarized in the following

1) The rising time for the hole size 25 mm is about 1.29 times longer than the scenario of hole size 90 mm. on the other hand the fountain height for larger hole size is approximately 2.1 times higher than the smaller hole size. Therefore, the hole size has a huge impact on the release rate of subsea gas, which increases with the size of the hole.

2) Water depth has a significant role on subsea gas release behaviors, this is obvious from our simulation results that the time needed for the gas plume reached sea surface when the water depth 13 m is about 3 times longer than the water depth 4 m whilst the fountain height for lower water depth is near to 1.2 times higher than the deeper water depth.

The scope of this study is modeling of underwater gas release with varying the effected parameters on underwater gas plume. So the future work is planned to build 3D CFD model with considering the effect of ocean current on the underwater gas plume.

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