Controlling Single Molecule Conductance by a Locally Induced Chemical Reaction on Individual Thiophene Units

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Abstract: Among the prerequisites for the progress of single-molecule-based electronic devices are a better understanding of the electronic properties at the individual molecular level and the development of methods to tune the charge transport through molecular junctions. Scanning tunneling microscopy (STM) is an ideal tool not only for the characterization, but also for the manipulation of single atoms and molecules on surfaces. The conductance through a single molecule can be measured by contacting the molecule with atomic precision and forming a molecular bridge between the metallic STM tip electrode and the metallic surface electrode. The parameters affecting the conductance are mainly related to their electronic structure and to the coupling to the metallic electrodes. Here, the experimental and theoretical analyses are focused on single tetracenothiophene molecules and demonstrate that an in situ-induced direct desulfurization reaction of the thiophene moiety strongly improves the molecular anchoring by forming covalent bonds between molecular carbon and copper surface atoms. This bond formation leads to an increase of the conductance by about 50% compared to the initial state.

One of the challenges associated with reducing electronic circuits to single-molecule components is the formation of reliable, low-resistance contacts between the molecule and the metallic electrodes to enable an efficient charge transport across the metal–organic interface.[1–3] While π-conjugated systems with a delocalized electronic structure are typically characterized by a high electronic conductance through the molecules,[4–6] they are usually equipped with non-conjugated end groups designed to selectively couple the molecules to specific metallic electrodes.[5–7] Moreover, the ability to control the junction’s geometry with the spatial overlap of orbitals is an important prerequisite for the optimization of the electronic device properties, such as the charge-injection efficiency and the charge mobility.[6–11] To this end, covalent bonds between the molecule and the metal electrodes are advantageous to ensure a robust mechanical and electronic connection. In the ideal case, the organo–metallic contacts of the molecular electronic element are established through a chemical reaction between the molecule and the metal. The

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most popular examples are systems with covalent bonds formed between sulfur-containing molecular end groups and gold electrodes.[14] When the anchoring is realized through direct covalent bonds between the carbon atoms and the metal, a higher electronic conductance is obtained due to a stronger coupling between the metal and the molecular \( \pi \)-system and a greater overlap of metal and molecular wave functions. The most common methods used to measure the electronic properties of organic molecules are based on the formation of break junctions[13] by trapping the molecules between contacts. This technique can be realized in a mechanically controlled break junctions, and in a more controlled way in scanning tunneling microscopes (STM)[14] or in conductive atomic force microscopes (AFM).[15] Variations of the molecular geometry and the number of molecules in the junction as well as the significant effects of the local environment may lead to very different signatures in the conductance even for the same molecular species. To avoid these uncertainties, we use an STM operating under ultra-high vacuum conditions and at low temperature. The instrument is particularly suitable for performing accurate conductance measurements through single molecules, as it allows us to characterize the structure at the submolecular level[16–19] and to atomically precisely contact one specific end of the molecular unit with the tip apex.[20–25]

Here we explore a new class of covalent bond formation between a chemically reacted thiophene molecular end group and Cu electrodes. We investigate the molecule-metal-electrode anchoring and the conductance across a tetraceno-thiophene (TCT) molecule, a pentacene analogue, by controlling the formation of covalent C–Cu bonds between the molecule and the metal substrate on the submolecular scale.

We demonstrate in this work that the resulting tetraceno derivative (TC-D) shows a significant increase in the conductance, as a result of a better contact formation between the bifurcated end of the remaining thiophene part and the metal surface, whose fingerprint is an intramolecular enhancement of the local density of states close to the Fermi energy.

Previously we have shown[26] an in situ induced direct desulfurization reaction on single thiophene units on Cu(111) surfaces driven by the electric field confined in the tunnel junction (Figure 1). This desulfurization reaction of the TCT molecules can also be induced thermally (see Supporting Information 2 for details). We now included high-resolution non-contact AFM (nc-AFM) images clearly revealing the opening of the thiophene ring and the split-off of the S-atom. Using a CO-functionalized tip for non-contact AFM, we are sensitive to the potential energy basins produced by the Pauli repulsion between the CO-terminated tip apex and specific atomic positions within adsorbed molecules.[27] This methodology gives rise to an extremely high resolution of intermolecular structures.[28] The high-resolution nc-AFM images of TCT and TC-D as well as simulated nc-AFM images obtained from DFT calculations clearly uncover the inter-molecular structure before and after the chemical surface reaction (Figure 1).

As a consequence of this reaction, the two carbon atoms formerly connected to the S atom in the thiophene group bind covalently to Cu surface atoms, while the S atom is irreversibly detached from the molecule. The resulting molecule, a tetraceno derivative (TC-D) is therefore anchored to the surface by two covalent bonds.

To probe the effect of the strong anchoring of one of the molecule’s end groups to the metal electrode, we performed conductance measurements through single TCT and TC-D molecules. The geometry of the measurement is sketched schematically in Figure 2. Before each conductance measurement, the molecule is first imaged by STM. Then, the tip is placed on top of the acene termination of the molecule, the feedback loop is switched off and the tip is moved towards the molecule while the current is recorded as a function of the tip-sample displacement (\( z \)). At the jump-to-contact position we obtain the conductance in units of the conductance quantum \( G_0 = 2e^2/h = 77.5 \, \mu\text{S} \). Figure 2a shows an example of such a \( G(z) \) curve. The exponential increase of the conductance corresponding to the tunneling regime is due to the reduction of the width of the tunneling barrier.

At a certain distance of the tip above the molecule, an abrupt increase in conductance is observed with a step-like...
The error bar corresponds to the standard deviation extracted from the system. In order to correlate the changes of the conductance to the intrinsic molecular properties, we calculated the projected density of states (PDOS) by summing over the atoms on different parts of the molecules (Figures 3b and c). An overall feature (Figure 2a). This corresponds to the jump-to-contact of the molecular moiety to the tip apex and indicates lifting-up of this part of the molecule from the surface. We speculate that the contact of the terminal acene ring to the tip is formed through the molecular π-orbitals. The conductance through the molecular system \((G_M)\) is extracted from the step height in the jump-to-contact regime (Figure 2a). Each conductance measurement was performed on a different molecule. The jump-to-contact behavior can occur at different tip-sample distances. To reduce the influence of the geometry of the junction with the molecule, the molecule presented from the surface on the conductance comparison, we only considered data obtained within a small tip-distance window of 50 pm (see Supporting Information 4 for details). All these conductance values \(G_M\) for TCT and TC-D are summarized in the histograms in Figure 2b. Since TCT is a pentacene analogue, we also added as a reference the values measured for pure pentacene. The average conductance values for TC-D are significantly larger; they are more than 45% higher compared with TCT and 130% compared to pentacene. The broad conductance distributions reflect the variation in the bonding geometries at the contact with the tip and different tip geometries encountered in different experiments.

We attribute the differences between the conductance through different types of molecules (TCT, TC-D, pentacene) to variations of the molecule-metal interaction strength together with changes in the molecular electronic structure. These effects are further analyzed using DFT calculations, in particular by comparing the electronic structure and the degree of the hybridization of the molecules to the Cu surface atoms (Figure 3). The conduction of the molecular junction is mainly determined by the electronic structure close to the Fermi level and by the transmission eigenchannels through the calculated molecular structures (ensuring a low molecular resistance).

In Figure 3a we show the partial charge densities of each type of molecule, as defined in the figure caption by energy integration around the Fermi level. For all three molecules the spatial distribution of the electronic charge shown in Figure 3a is rather uniform along the tetracene part. However, there is a significant difference on the reacted thiophene side of the TC-D molecule. The locally driven desulfurization reaction leads to the formation of covalent bonds between the bifurcated termination and the copper surface. The covalent bond formation results in a strong increase in the overlap of the wave functions between the metal electrode and the molecular end group, thus inducing a strong anchoring process that may also induce a reduction of the contact resistance (for details see Supporting Information 5). The strong coupling between the molecular end group and the surface atoms is also evidenced in the experiments by the fact that the reacted molecule cannot be picked up completely with the STM tip apex by STM manipulation, whereas this is possible for TCT and pentacene molecules.

In order to correlate the changes of the conductance to the intrinsic molecular properties, we calculated the projected density of states (PDOS) by summing over the atoms on different parts of the molecules (Figures 3b and c). An overall enhancement of the PDOS close to the Fermi level \((E_F)\) of the metal is observed on the entire TC-D molecule, which is related to the electronic-density delocalization of the aromatic system, but it is most significant at the thiophene side (side 2) compared to pentacene and TCT. The enhancement of the PDOS close to \(E_F\) for TC-D is responsible for the...
which was cleaned by repeated cycles of Ar\(^+\) high-vacuum (UHV) conditions. The substrate is a Cu(111) surface, methodology described in reference \[29\]. The spectroscopic investigations of the identity and purity of tetracenothiophene and the precursors is presented in the Supporting Information 1. The pentacene and TCT molecules were thermally sublimated separately and in different experiments from ceramic crucibles at 445 K and 485 K, respectively. The molecular coverage in different experiments was between 7 to 15\% of a monolayer. We have always selected isolated molecules for the conductance measurements which is one of the advantages using STM. For the pentacene sublimation, the substrate was at room temperature, while for TCT the Cu(111) surface was held at 200–250 K during deposition to avoid the thermally induced desulfurization (see Supporting Information 2). After each evaporation, the sample was transferred in situ into the STM, operating at low temperatures (6 K). The STM measurements were carried out by applying the bias voltages to the sample. The TC-D molecules were produced in the STM by the direct desulfurization reaction, performed by placing the tip on top of the thiophene moiety of a TCT molecule, at a height corresponding to the set point of 100 pA, 100 mV. Then, the feedback loop was switched off and at constant height the voltage was ramped. Changes in the \(I(V)\) curves were observed when the reaction occurred. The high-resolution nc-AFM frequency-shift images were acquired with a CO-terminated tip. The CO molecule was vertically manipulated to the tip apex from the NaCl(100) islands. For the AFM image acquisition, the tip was positioned over the thiophene group in the case of the TCT molecule and over the tetracene termination in the case of the TC-D molecule. At the feedback 150 pA and 10 mV, the STM Z-feedback was switched off and the bias voltage was reduced to 1 mV. The images were obtained with both amplitude and phase feedback of the AFM enabled and using the oscillation amplitude of 40 pm. The conductance measurements were performed by a vertical manipulation procedure. After switching off the feedback loop at 100 pA, 100 mV and moving the tip toward the molecules by 250 pm, the \(I(z)\) curves were recorded above the acene-terminated end of each molecule, which were identified beforehand locally by imaging the molecule in the STM mode. STM images of a TCT molecule and of the produced TC-D molecules are shown in Supporting Information 3.

**Computational Methods:** DFT-vdW calculations were performed using the Vienna Ab-initio Simulation Package (VASP)\(^{[30,31]}\) An (8 \times 6) periodic supercell with a 4-layer slab was employed to model the adsorption of each type of molecule (TCT, TC-D, pentacene) on the Cu(111) surface. For each geometry optimization, full relaxation of all degrees of freedom of the molecule’s atoms was performed, and the projected densities of states (PDOS) over the p\(_z\) orbitals of the carbon atoms were computed for the equilibrium geometries. Ion-electron interactions were described with the projector augmented-wave (PAW) method,\(^{[32]}\) and the exchange correlation was modelled within the generalized gradient approximation (GGA).\(^{[33]}\) Van-der-Waals dispersion forces were included using the vdW-DF-cx method.\(^{[34,35]}\) It is worth mentioning that with this functional we obtained a 2.38 Å adsorption distance of pentacene on Cu(111), in good agreement with experimental findings of 2.34 ± 0.02 Å.\(^{[36]}\) We

![Figure 3. Calculated electronic structure and adsorption conformation of pentacene, TCT and TC-D molecules on Cu(111).](image)

The high-resolution nc-AFM frequency-shift images were acquired with a CO-terminated tip. The CO molecule was vertically manipulated to the tip apex from the NaCl(100) islands. For the AFM image acquisition, the tip was positioned over the thiophene group in the case of the TCT molecule and over the tetracene termination in the case of the TC-D molecule. At the feedback 150 pA and 10 mV, the STM Z-feedback was switched off and the bias voltage was reduced to 1 mV. The images were obtained with both amplitude and phase feedback of the AFM enabled and using the oscillation amplitude of 40 pm. The conductance measurements were performed by a vertical manipulation procedure. After switching off the feedback loop at 100 pA, 100 mV and moving the tip toward the molecules by 250 pm, the \(I(z)\) curves were recorded above the acene-terminated end of each molecule, which were identified beforehand locally by imaging the molecule in the STM mode. STM images of a TCT molecule and of the produced TC-D molecules are shown in Supporting Information 3.

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considered a 500 eV energy cutoff in the plane-wave expansion, with a 2 × 3 k-point mesh in the 1 × 1 unit cell as sampling of the Brillouin zone reciprocal space. The electronic convergence criterion was 1 × 10⁻⁵ for all static calculations, and the convergence on forces in the relaxations was 0.05 eV/Å⁻¹. The VESTA software[30] was employed for the ball-and-stick and charge density representations.

From the DFT-vdW optimized adsorptions geometries, AFM image simulations were performed using the probe-particle model,[32,36] which is based on the van-der-Waals (vdW) and electrostatic interactions between the sample and the tip. The electrostatic interaction between tip and sample was calculated from the surface Hartree potential obtained from DFT calculations, while the classic pairwise potential (Lenard–Jones) was used to describe the Van-der-Waals attraction and the Pauli repulsion. The AFM calculations were carried out with different values of the effective charge of the probe particle in order to obtain AFM images that show a good agreement with the experimental results. It was found that a quadrupole (d₂) distribution of the effective charge was more appropriate to describe the CO-terminated tip.[37] Note that all theoretical simulations were performed with a fixed lateral stiffness of (k₀) = 0.25 Nm⁻¹ and by approximating the CO-functionalized tip by an effective atomic radius R₀ = 1.661 Å. The agreement between the theoretical findings and the experiments enables us to understand and interpret the origin of the chemical contrast.

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Conflict of interest

The authors declare no conflict of interest.

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