Probing Neural Networks for Dynamic Switches of Communication Pathways

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Abstract

Dynamic communication and routing play important roles in the human brain to facilitate flexibility in task solving and thought processes. Here, we present a network perturbation methodology that allows to investigate dynamic switching between different network pathways based on phase offsets between two external oscillatory drivers. We apply this method in a computational model of the human connectome with delay-coupled neural masses. To analyze dynamic switching of pathways, we define four new metrics that measure dynamic network response properties for pairs of stimulated nodes. Evaluating these metrics for all network pathways, we found a broad spectrum of pathways with distinct dynamic properties and switching behaviors. Specifically, we found that 60.1\% of node pairs can switch their communication from one pathway to another depending on their phase offsets. This indicates that phase offsets and coupling delays play an important computational role for the dynamic switching between communication pathways in the brain.

Author summary

A big challenge in elucidating information processing in the brain is to understand the neural mechanisms that dynamically organize the communication between different brain regions in a flexible and task-dependent manner. In this theoretical study, we present an approach to investigate the routing and gating of information flow along different pathways from one region to another. We show that stimulation of the brain at two sites with different frequencies and oscillatory phases can reveal the underlying effective connectivity. This yields new insights into the underlying processes that govern dynamic switches in the communication pathways between remote sites of the brain.

Introduction

Over the past decades it has been shown that the brain, facing a specific task or not, reveals a well-structured functional organization \cite{1,2}. This has been specifically...
investigated for resting-state networks [4–7], but also for other networks when the brain is performing different tasks [6,8]. These findings lead to the idea that resting-state networks describe an inherent functional organization of the brain which is optimized to perform a wide range of tasks it encounters frequently [9,10]. If faced with a task that requires synchronization between brain areas not typically coupled at rest, this organization has to be altered temporarily in order to perform that task efficiently [11,12].

Within a complex network like the human brain, multiple structural pathways exist between most pairs of nodes given a sufficiently high spatial resolution. Since synchronization along such pathways seems to play an important role in the formation of functional clusters, we set out to identify general principles of how these pathways interact with each other during synchronization processes. Several studies have emphasized the importance of information transmission delay for synchronization processes as well as its role in the formation of functional clusters in the brain [3,13–19]. Thus, we hypothesize the time lag inherent to a communication path to be a key factor in the interaction between multiple pathways. These time windows are determined by axonal signal transmission delays as well as rise and decay properties of the post-synaptic response. Here, we focus on the former, expecting region-specific differences in the latter to be negligible for the long-range connections considered in this study. In particular, we predict that two brain regions trying to communicate at a certain frequency with a given phase offset will use only a fraction of their available communication paths. Further, we predict that the selection of communication paths will be influenced by their interaction time windows.

To test these hypotheses, we approximate cortico-cortical communication processes via a computational network model of delay-coupled, oscillating nodes. We introduce an extrinsic stimulation set-up for this model that allows to detect network interactions between pairs of nodes. This stimulation approach relies on the entrainment of a given pair of nodes to oscillate at the same frequency, but with a certain phase lag relative to each other. Comparing the coherence along different pathways over different stimulation phase offsets then reveals the phase preferences for different routes. While Fig 1A illustrates the extrinsic stimulation setup, Fig 1B motivates the use of different stimulation phase lags. It is important to note that even in the absence of any interaction through the network, there might be some induced trivial coherence between two stimulated nodes due to the external signal (Fig 1B). Thus, the coherence is measured for many different stimulation phase offsets and the measurement with the lowest coherence is chosen as the baseline. Any deviation in the coherence from this baseline can be attributed to induced changes in the coupling between the two stimulated nodes through the network, which may happen due to a switching in the pathways (compare Fig 1C and 1D). We propose that these differences in phase preferences at different pathways act as a switching and gating mechanism used by the brain to establish communication between remote brain areas when needed. Our method allows to investigate these mechanisms by probing the network for these dynamic switches in communication pathways.
Fig 1. Illustration of the stimulation methodology and different possible outcomes. (A) The stimulation at the two nodes could potentially activate different paths (blue or red) in the network. (B) As an alternative explanation it could also be possible that a coherence between the two nodes is induced even in the absence of any direct interaction in the network. (C) In this example a stimulation phase offset of \(0\pi\) induces a network interaction between the two stimulated nodes through path 1 (blue). (D) In contrast, when stimulating with a phase offset of \(\pi\) another path is activated (red).

Computational Model

Our computational model (Fig 2) is based on the widely used Jansen-Rit neural mass model [20] which employs a mean-field approach to model the interaction between cell populations in the infragranular (depicted in green), granular (depicted in blue), and supragranular (depicted in red) layer, as illustrated with the relevant equations in the Figure. The function \(\sigma(V)\) that transforms average membrane potentials to firing rates is a parameterized sigmoid (depicted in cyan). The standard parametrization originally proposed by Jansen and Rit reflects cortical oscillatory activity in the alpha frequency band. These parameters were chosen based on experimental findings in the neuroscience literature and are reported in Table 2 in the Methods and Materials section. Since the purpose of this article is the investigation of the effect of pathway time-scales on neural synchronization processes and not the effect of node time-scales, we decided to use this standard parametrization for each node in our network [20].

For the purpose of investigating networks of extrinsically perturbed Jansen-Rit models, the following two extensions were added: First, we coupled multiple Jansen-Rit nodes via delayed, weighted connections between their infragranular pyramidal cell populations (depicted in yellow). Secondly, weak external drivers were applied at two stimulation sites influencing the average membrane potential of the infragranular layer with phase offset \(\Delta\varphi\) between the two drivers (depicted in purple).

Results

We first present results of simulations in simple networks of only 2 or 3 nodes. Specifically, we show how the coherence between a pair of stimulated nodes depends on the transmission delay of their connections and on the relative phase offsets between the two stimulation signals. Subsequently, we move on to a model of the cortico-cortical synchronization processes within a single hemisphere of the human brain. Cortico-cortical coupling strengths and delays were informed by an approximation of the human connectome obtained from diffusion tensor imaging (DTI) data. Additionally, we fit the model to match functional connectivities calculated from electroencephalography (EEG) recordings. In this model, we show how the coherence
Inhibitory cells in supragranular layers
\[ \frac{dV_{i3}}{dt} = I_{j3} \text{ with } \frac{dI_{j3}}{dt} = -\frac{2}{\tau_e} V_{j3} - \frac{H_e}{\tau_e} S_{j3} \]
→ firing rate \( S_{j3} = \sigma(V_{j3}) \)

Excitatory spiny cells in granular layers
\[ \frac{dV_{i2}}{dt} = I_{j2} \text{ with } \frac{dI_{j2}}{dt} = -\frac{2}{\tau_e} V_{j2} - \frac{H_e}{\tau_e} S_{j3} \]
→ firing rate \( S_{j2} = \sigma(V_{j2}) \)

Excitatory pyramidal cells in infragranular layers
\[ \frac{dV_{i3}}{dt} = I_{j3} \text{ with } \frac{dI_{j4}}{dt} = -\frac{2}{\tau_e} V_{j4} - \frac{H_e}{\tau_e} \left( S_{j2} + \epsilon_{ext} I_{j4} \right) \]
\[ \frac{dV_{i4}}{dt} = I_{j4} \text{ with } \frac{dI_{j4}}{dt} = -\frac{2}{\tau_e} V_{j4} - \frac{H_e}{\tau_e} \left( S_{j2} + \epsilon_{ext} I_{j4} \right) \]
→ firing rate \( S_{j3} = \sigma(V_{j3} - V_{j4} + \epsilon_{ext} V_{j4}) \)

Activation Function:
\[ \sigma(V) = \frac{2e_0}{1 + e^{-\sigma(V_0-V)}} \]

External perturbation
If stimulus at node \( j \) or 0 otherwise

Connectome interactions
\[ S_{j,\text{in}}(t) = \sum_{m=1}^{N} C_{m,j} S_{m3} \left( t - \frac{L_{m,j}}{v} \right) \]
with conduction delays \( L_{m,j}/v \)

Fig 2. Neural mass model with external stimulation. A schematic of the neural mass model showing the interactions between the three neuronal populations in the infragranular, granular, and supragranular layer. Each post-synaptic potential is modeled using two differential equations (average membrane potentials \( V \) and average synaptic currents \( I \)). Several of these neural mass models are interacting through a connectivity matrix (yellow). The external perturbation (purple) modulates the average membrane potential of pyramidal cells at 2 nodes in the network.

between stimulated nodes changes over phase lags and how this effect relates to the connectedness and distance of the nodes. In a final step, we identify the pathways responsible for the interaction between the stimulated nodes, analyze their phase lag preferences and identify cases of phase-related switching between pathways. For all simulations, we use the computational model defined in the previous section.

Simple Models With 2 or 3 Nodes

The idea behind the extrinsic stimulation approach can be well explained using a simple toy-model of 2 directly coupled Jansen-Rit nodes, where each node is stimulated with a \( f_{\text{ext}} = 11 \) Hz sinusoidal signal with strength \( c_{\text{ext}} = 0.25 \) mV. Fig 3 shows the coherence between the driven nodes for systematic changes in the phase offset between the stimuli and the distance between the coupled nodes. While uni-directionally coupled nodes can have preferences for any stimulation phase offset, as shown in Fig 3A, bi-directionally coupled nodes are more susceptible for stimulation at in- or anti-phase (see Fig 3B). This shows that the communication between coupled pairs of nodes can be modulated by stimulation and that communication channels can have characteristic stimulation phase offset preferences, depending on their length [21].

To quantify the modulation of communication, we define the pathway-synchronization-facilitation (PSF), measuring for a given pair of weakly stimulated network nodes \( k_i \) and \( k_j \) how their interaction is dependent on specific phase offsets:

\[
\text{PSF}(k_i, k_j) = \max_{0 \leq \Delta \phi \leq 2\pi} \left( \text{coh}(k_i, k_j, \Delta \phi) \right) \quad (1)
- \min_{0 \leq \Delta \phi \leq 2\pi} \left( \text{coh}(k_i, k_j, \Delta \phi) \right), \quad (2)
\]

where \( \text{coh}(k_i, k_j, \Delta \phi) \) is the coherence between network nodes \( k_i \) and \( k_j \) for stimulation phase offset \( \Delta \phi \). The PSF is high for node pairs if their coherence is high.
**Fig 3.** Simulations with 2 or 3 nodes, where the red edges correspond to the distance that was varied. Color indicates the coherence between the two driven nodes. The PSF values are shown at the bottom of each panel. (A) Nodes with direct uni-directional coupling. (B) Nodes with direct bi-directional coupling. (C) Nodes with indirect (via a third node) bi-directional couplings. The intermediate node was placed at 25% of the total connection distance, while the overall distance between the outer nodes was varied. (D) Nodes with indirect bi-directional couplings where the overall distance was kept at 100 mm, while the intermediate node was positioned at varying positions along the connection. Parameters used in all panels: $v = 3 \text{ m/s}$, $C_{m,j} \cdot c_{\text{net}} = 0.1$ if there is a connection from node m to j or 0 otherwise.

for one stimulation phase offset and low for another, i.e., the relative phase of the stimulation at the two sites matters strongly. The PSF curves in Fig 3A and 3B show that in both cases there is a PSF effect ($\text{PSF} > 0$) and in the case of bi-directionally coupled nodes the strength of this effect depends on the distance between the nodes.

To extend this idea to communication via indirect pathways, we investigated synchronization between 2 nodes connected only indirectly via a third intermediate node. We used bi-directional couplings for both connections and both end nodes were stimulated as described previously. As can be seen in Fig 3C and 3D, the interaction between the two weakly stimulated nodes not only depended on the length of the connection, but potentially also on the relative position of the third node on the indirect path.
Fig 4. Pairwise measures of connectivity and distance between all 33x33 region pairs. (A) Structural connectivity matrix with all connections smaller than 0.1 set to 0. (B) Inter-regional distances in mm. (C) Functional connectivity matrix derived from coherence of EEG data bandpass-filtered around 10 Hz. (D) Functional connectivity derived from coherence of neural mass model simulations bandpass-filtered around 10 Hz.

Connectome Model Without Stimulation

As shown in the previous section, the coherence in a network of only three nodes can already exhibit very complex dependencies on the stimulation phase offset. Next, we wanted to analyze network communication patterns in the case of a complex network with multiple competing pathways. To this end, we used a model of 33 delay-connected nodes, representing one hemisphere of the human connectome [19]. The structural connectivity matrix was obtained from DTI-based tractography data as described in more detail in the Methods and Materials section. Fig 4A shows the sparse connectivity matrix $C_{mj}$ used to connect the 33 regions and Fig 4B the corresponding distances $L_{mj}$ between region pairs. These two matrices are used for connecting the different populations of excitatory pyramidal cells (green in Fig 2) in a delay-coupled network of a single-hemisphere.

For the same 33 regions, EEG resting-state recordings from the same subjects were used to calculate pairwise coherences in the 10 Hz range as shown in Fig 4C. Similarly, we simulated the 33 connected neural-mass-models and processed the time-series of the pyramidal cells in the same way as the EEG data. This yielded a 33 x 33 functional connectivity matrix which we compared to the empirical functional connectivity by calculating the Pearson correlation coefficient.

The selection of parameters was based on the rationale to match the functional connectivity observed in the network model as good as possible to empirical EEG-based functional connectivity from human subjects. We performed a grid search over global structural connectivity scaling $c_{net}$ and transmission velocity $v$ to obtain the best match between modeled and empirical data. By fitting the velocity, we ensured that our pathway delays reflect realistic, empirically observed timescales of cortico-cortical interactions. We found the highest correlation ($r = 0.57$) for $c_{net} = 20$ and $v = 3$ m/s, so that we used these parameters for subsequent analyses (Fig 4D). Notice that this
Fig 5. Stimulation parameter evaluation. (A) Coherence between the stimulus and the stimulated region for varying stimulus strength and frequency. The color corresponds to a mean value over 5 simulations at randomly chosen network nodes. (B) Coherence between the stimulus and the full network averaged over all 33 nodes for varying stimulus strength and frequency. The coherence was calculated between the stimulus and each region and then averaged. (C and D) Coherence between two stimulated nodes for varying stimulation strength and phase offset. C and D correspond to two different node pairs that were stimulated demonstrating the variability in phase lag preferences expressed by different pairs.

correlation is comparable with values of other bottom-up models reported in the literature [19,22], which is remarkable considering that we set a substantial amount of structural connections to 0. Interestingly, both simulated and empirical functional connectivity correlate stronger with the inter-regional fiber distances ($r_{sim} = -0.65$, $r_{emp} = -0.70$) than with the structural connectivity ($r_{sim} = 0.65$, $r_{emp} = 0.51$), highlighting the importance of signal transmission delays for cortical coupling processes.

Connectome Model With Stimulation

Based on this model of cortical activity we used the stimulation approach to investigate how pathways facilitate synchronization between network nodes at certain phase lags between the nodes. Specifically, we weakly stimulated different pairs of cortical regions with varying phase offsets between the two stimulation signals while measuring the coherence between the stimulated nodes at each phase offset. As argued above, finding differences in the coherence over stimulation phase offsets would indicate phase-specific communication modulation between the stimulated nodes. Before analyzing PSF effects in the connectome model, it was necessary to determine the optimal stimulation frequency and strength for this model. This was performed in two steps. First, we stimulated a single region in our network with a stimulus of varying frequency (4-22 Hz) and strength (0.01-2 mV) while evaluating the coherence between region and stimulus. The mean coherence (mean over 5 different stimulated nodes) for each parameter combination can be observed in Fig 5A. They reflect the well-studied relationship between a driver and an oscillator described by the so-called Arnold tongue [23]. Since our main analysis will focus on coupling effects through different network paths between...
two stimulated nodes, we also calculated the coherence between stimulus and all network nodes (Fig 5B). This average coherence to the full network was strongest at 9-11 Hz, which is also the intrinsic frequency of unperturbed network nodes [24]. Interestingly, at this frequency the coherence to the directly stimulated node was weakest (compare Fig 5A). Based on this, we set the frequency of our stimulus to 11 Hz, at which the network (and not only the directly stimulated node) was most susceptible for entrainment by an external stimulation.

In a second step, we stimulated pairs of nodes with 11 Hz stimuli. We varied the stimulus strength (0.25 - 1 mV) and the relative phase offset between the stimuli (0 − 2π), while evaluating the coherence between the stimulated nodes. All other parameters were chosen to be the same as for the previous simulation. As can be seen in Fig 5C and Fig 5D, the variance of the coherence over phase offsets depended on the stimulation strength. Based on visual inspection of the coherence patterns of 20 different region pairs, we chose our stimulus scaling to be \( c_{ext} \) = 0.5 mV, leaving the variance of the post-synaptic potential of the neural masses in a biologically plausible range and such that the external driver is relatively weak in comparison to the internal network dynamics. This gave us the final set of global model parameters which were used throughout all subsequent simulations and are reported in Table 1.

### Table 1. Model Parameters

| Param. | Value | Interpretation |
|--------|-------|----------------|
| \( c_{net} \) | 20 | global connection strength scaling |
| \( v \) | 3 m/s | global velocity scaling |
| \( f_{ext} \) | 11 Hz | extrinsic stimulation frequency |
| \( c_{ext} \) | 500 \( \mu \)V | amplitude of extrinsic stimuli |

The variability in the coherence between stimulated region pairs that we observed not only over stimulation phase offsets but also over different pairs (as depicted for 2 exemplary region pairs in Fig 5C and 5D) shows that the stimulated region pairs interacted with each other and that the interactions showed a characteristic profile of phase offset preferences. To statistically confirm the variance in the coherence between stimulated region pairs over phase offsets, we ran simulations with stimulations of each possible node pair. Again, we varied the phase offset between the two stimuli (16 equally spaced phase offsets between 0 and 2π) and evaluated the coherence between the stimulated nodes for each phase offset. Subsequently, those coherences were used to calculate the PSF for each region pair as defined in Eq (1). Using a one-sample t-test, we found the PSF effect to be significantly larger than zero (mean = 0.1567, CI = [0.1445,0.1689], \( t = 25.2595, p < 0.0001 \)). Hence, we were able to show with our extrinsic stimulation approach that pathways facilitated synchronization between cortical nodes and that the facilitatory strength depended on the phase lag between the region’s average PSPs.

With the PSF effect established, we continued by investigating its dependence on certain features of the underlying structural connectivity graph. For this purpose, we searched for all possible pathways between each pair of stimulated nodes based on the structural connectivity matrix reported in Fig 4A. Since every stimulated pair of nodes was connected by at least one path via at most 6 edges, we restricted the search to pathways including 6 edges at maximum. With these pathways at hand, we started out by evaluating how the PSF effect changed with increasing network distance. An analysis of variance showed that the effect of shortest path length (minimum number of edges separating a pair) on log(PSF) was significant, \( F(5,521) = 97.6141, p < .0001 \). As can be seen in Fig 6A, we observed the trend that the PSF effect decreases with the number of nodes separating the stimulated nodes. Furthermore, as depicted in Fig 6B, this trend was supported by a significant correlation between the PSF effect and the
Fig 6. Pathway Synchronization Facilitation (PSF). All panels show the PSF on a logarithmic vertical axis and red areas indicate 95% confidence intervals. Blue boxes indicate 1 standard deviation. (A) Dependence of PSF effect on shortest path length between the node pairs. (B) Correlation between PSF effect and average fiber length along the shortest pathway between the stimulated nodes. (C) Dependence of the PSF effect on the number of paths connecting the stimulated nodes.

length of the shortest pathway between the stimulated nodes ($r = -0.56, p < .0001$), a measurement that is strongly related to both interregional distance and minimal number of separating edges. Thus, we conclude that there is a tendency for a decrease in the interaction of stimulated node pairs with increasing network distance between the nodes, where network distance can be measured either as the number of edges or as the summed up length of the edges of the shortest pathway connecting the nodes.

Next, we investigated the dependence of the PSF effect on the connectedness between the stimulated nodes. An analysis of variance showed that the effect of the number of connecting paths (only counting paths with 5 edges or less, all nodes with more than 5 connecting paths were pooled into one level) on log(PSF) was significant, $F(5,501) = 10.0827, p < .0001$. The latter result can be observed in detail in Fig 6C.

Evaluation of Pathway Activation

Having described the influence of the external driver on the coherence between stimulated nodes, we next identified which particular pathways were involved in this interaction. For this analysis, we define the pathway activation (PA) for a pathway through $n$ nodes $k_i$ with $i = 1..n$ at a phase offset $\Delta \varphi$ as the minimum of the pairwise coherences between neighboring pathway nodes:

$$PA(k_1..k_n, \Delta \varphi) = \min_{i=1,...,n-1} \left(\text{coh}(k_i, k_{i+1}, \Delta \varphi)\right).$$

(3)

In other words, if communication fails at any point along a pathway, leading to a reduced coherence between the involved nodes, this is considered to be a bottleneck for the information flowing through that pathway. We evaluated the pathway activation (PA) for all pathways of up to $n = 5$ nodes connecting a given pair of stimulated nodes for all stimulation phase offsets. Doing this for each stimulated node pair, we found different classes of pathway interactions: Some pairs show only a very small selectivity for the stimulation phase offset (Fig 7A), while other node pairs were connected by paths with PA values with a strong dependence on the phase offset (Fig 7B, 7C). Moreover, some of these node pairs switched their interaction between different pathways depending on the stimulation phase offset, as shown in Fig 7C and the two switching pathways in Fig 7D and 7E.

To further analyze how the communication via specific pathways depends on the stimulation phase offset, we define the pathway phase selectivity (PPS) of a pathway $P_1$...
Fig 7. Dependence of path activation (PA) on stimulation phase offset. Panels (A–C) show the PA values (radius) for different stimulation phase offsets (angle) for exemplary node pairs. Blue corresponds to the pathway with the strongest overall PA. Red corresponds to the second strongest pathway that has no overlapping segments with the strongest. Green curves show the strongest of the remaining pathways (with possible overlapping path segments with the former two). The two arrows in (C) indicate the phase offsets which are used in panels (D–E). (D) Connectome pathways for stimulation of the node pair shown in (C) at phase offset $1.5\pi$. The two stimulated nodes are shown as two black dots. The most active pathway at this stimulation phase offset is highlighted with a stronger line width. All colors correspond to the coherence of nearest neighbours in the connection graph. (E) Similar to (D) but for stimulation phase offset $0.5\pi$ with a different most active pathway. (F) Histogram of pathway phase selectivity. The values of the strongest paths of examples shown in (A–C) are marked with arrows. (G) Histogram of phase differences between the stimulation phase offset where the most active pathway has the highest pathway activation and the stimulation phase offset where the coherence between the stimulated nodes was highest. (H) Similar to (G) but for the second most active pathway (excluding path overlaps with the most active pathway). (I) Histogram of normalized pathway switching index between the strongest and second strongest pathways. The values of the node pairs of examples shown in (A–C) are marked with arrows. The square-root normalization of the PSI transforms back from the space of multiplied coherence values to the original non-squared coherence space (analogous to a transformation from variance to standard-deviation).
similar to the PSF as

\[ PPS(P_1) = \max_{\Delta \varphi} (PA(P_1, \Delta \varphi)) - \min_{\Delta \varphi} (PA(P_1, \Delta \varphi)), \]  

(4)

Pathways with relatively constant PA values for all stimulation phase offsets have a low PPS (example in Fig 7A), while pathways with a high variation in the PA values have a high PPS (example in Fig 7B and 7C). The evaluation of PPS values for all node pairs results in a bimodal distribution (Fig 7F). The activation of pathways in the first mode at \( PPS = 0.1 \) is very hard to influence with phase offsets. But we also found many node pairs with pathways in the second mode at \( PPS = 0.35 \). The communication of these later node pairs can be modulated using different phase offsets.

In a next step, we analyzed the relationship of pathway-specific phase preferences (as shown in Fig 7A-C) to the phase preferences of the stimulated nodes (as shown in Fig 5C-D). We chose the most active pathway per node pair (averaged over all stimulation phase offsets) and calculated the phase difference between the stimulation phase offset with the highest coherence and the stimulation phase offset with the highest PA. The histogram of these phase differences is significantly different from uniform, \( \chi^2(15, N = 514) = 273.05, p < .001, \) and has a peak at 0 (Fig 7G). In contrast, a similar analysis for the second strongest pathway (excluding all pathways with overlapping sections with the strongest path), results in a histogram that does not differ from a uniform distribution, \( \chi^2(15, N = 514) = 14.06, p = 0.52 \) (Fig 7H). Therefore, we conclude that the pathway with the strongest PA shows a similar phase preference as the coherence between the two stimulated nodes.

Finally, we quantified the switching between the strongest and second strongest pathway per node pair. To this end, we define the pathway switching index (PSI) between pathways \( P_1 \) and \( P_2 \) as

\[ PSI(P_1, P_2) = \max_{\Delta \varphi} (PA(P_1, \Delta \varphi) - PA(P_2, \Delta \varphi)) \]  

\[ \cdot \max_{\Delta \varphi} (PA(P_2, \Delta \varphi) - PA(P_1, \Delta \varphi)), \]  

(5)

(6)

The PSI is positive if the two pathways switch their activation depending on the stimulation phase offset, meaning that at one phase offset the first path is more active and at another phase offset the second path is more active. We found that 60.1% (309 of 514) of node pairs have a positive PSI between their non-overlapping strongest and second strongest pathways (Fig 7I). These results suggest that in this network of 33 nodes of the human connectome many node pairs have the capacity to switch their communication between at least two different pathways with a PA characteristic similar to the example shown in Fig 7C.

**Discussion**

We have carried out a computational study of cortico-cortical synchronization processes that strongly emphasizes the role of phase relationships for dynamic switches in communication pathways. In this process, we introduced a novel method to detect network interactions between pairs of cortical regions via an extrinsic stimulation scheme. Using our method, we were able to quantify the influence of different pathways on cortico-cortical synchronization processes between all pairs of 33 brain regions and could further identify the pathways those region pairs use to interact with each other. These pathways represent communication channels with distinct interaction time windows.

One of the main findings of our work is that we found the ability of regions to communicate via these channels to depend on the phase lag at which they try to
This finding is in line with the communication-through-coherence theory that predicts neural communication to critically depend on oscillatory phase differences and has received support from various experimental results [18]. In a transcranial magnetic stimulation study, Elswijk and colleagues demonstrated the effect of the stimulation on primary motor cortex to depend on the oscillatory phase of the latter [25]. Similarly, Helfrich and colleagues found the performance in a visual detection task to be modulated by the phase of a 10 Hz transcranial alternating current stimulation applied to the parieto-occipital cortex [26]. Together, these experimental results support the idea of the communication-through-coherence theory that the effectiveness of neural communication between a source and a target population is modulated by the oscillatory phase of the latter. In a study using multielectrode recordings in frontal eye field and area V4 of monkeys, Gregoriou and colleagues were able to show that neural populations at both recording sites synchronized at characteristic time lags when processing a stimulus in their joint receptive field [27].

Based on the communication-through-coherence theory this may be explained by the axonal and synaptic delays of the communication channel between these two areas, which render synchronization at this time lag most effective. The work presented in this article contributes to this line of arguing via its systematic investigation of the dependence between communication channel delays and neural synchronization. By finding that the synchronization of two brain regions via simultaneous extrinsic stimulation has a preference for the phase relationship between the two stimulation signals, we demonstrated the mechanistic role of phase lags for neural communication within a computational model. We were able to extend this finding to single communication channels, i.e. we showed that these channels can express more or less strong preferences for the phase lag between the brain regions they connect. On the one hand, this emphasizes the potential importance of communication channel delays for shaping the synchronization phase lag patterns of the brain. On the other hand, it suggests the effectiveness of information exchange via a certain communication channel in the brain to depend on the phase relationship between the communicating sites.

The functional significance of these phase lag preferences is underlined by our finding that different communication channels may be employed at different stimulation phase offsets between the communicating regions. We believe that the switching between different synchronization phase lags could be an effective mechanism through which brain regions can dynamically change their communication channels. Both external stimuli and network intrinsic signals could act as phase-resetting mechanisms at the communicating sites and thus switch from one communication channel to another within a few oscillatory cycles [28]. As suggested in [10,29], such a mechanism would provide the necessary flexibility to allow for the dynamic binding of remote neural representations into different combined representations.

Taken together, our results suggest a potential mechanism the human brain might have developed to use the physiological constraints imposed by coupling delays to its computational advantage. They motivate future brain stimulation studies to investigate phase-lagged neural synchronization, e.g., through multi-site transcranial stimulation, or optogenetics in combination with multielectrode recordings [30]. Thereby, the stimulation method we applied to our computational model could guide the development of brain stimulation protocols to probe dynamic switching between communication pathways in the brain. Additionally, our stimulation method and graph metrics are applicable in future theoretical studies characterizing the dynamic properties of network graphs.
Materials and Methods

In the following, we present the detailed parameters of the neural mass model, the preprocessing of the structural connectivity, and the functional connectivity.

Neural Mass Model

All simulation results reported refer to 16 minutes of simulated network behavior, using an explicit Euler method with an integration step-size of 0.5 ms. The parameters of the neural mass model are shown in Table 2.

| Param. | Value | Interpretation |
|--------|-------|----------------|
| $H_e$  | 3.25 mV | avg. gain of excitatory synapses |
| $H_i$  | 22 mV  | avg. gain of inhibitory synapses |
| $\tau_e$ | 10 ms | lumped time constant of excitatory synapses |
| $\tau_i$ | 20 ms | lumped time constant of inhibitory synapses |
| $c_1$ | 135 | avg. number of contacts from pyramidal cells to exc. interneurons |
| $c_2$ | 0.8 · $c_1$ | avg. number of contacts from exc. interneurons to pyramidal cells |
| $c_3$ | 0.25 · $c_1$ | avg. number of contacts from pyramidal cells to inh. interneurons |
| $c_4$ | 0.25 · $c_1$ | avg. number of contacts from inh. interneurons to pyramidal cells |
| $e_0$ | 2.5 Hz | maximum scaling of the synaptic gain |
| $r$ | 0.56 mV$^{-1}$ | steepness of the sigmoid function |
| $V_0$ | 6 mV | value with 50% of max. firing rate |
| $u_j$ | 120 - 320 Hz | sub-cortical noise distributed uniformly |

In the Jansen-Rit model, signal transmission between cell populations is realized via a convolution of average pre-synaptic firing rates with a post-synaptic response kernel. These convolution operations are mathematically described by the coupled ordinary differential equations in Fig 2 (each line describes a synaptic convolution operation) and turn pre-synaptic firing rates into post-synaptic potentials. The simple exponential form of the response kernel is in line with empirically measured post-synaptic responses and provides a sufficient approximation of synaptic response time scales for our purposes. To translate post-synaptic responses back into firing rates, an instantaneous sigmoidal transform is used as shown in Fig 2.

Structural Connectivity and Distance Estimates

In a first step to building a bottom-up model of cortical activity we needed to approximate the structural connections between different brain regions. As mentioned in the introduction, this can be done via DTI recordings. However, there are several technical limitations as to what extend human SC can be approximated based on DTI, and it depends on the systematic underestimation of inter-hemispheric connections. Thus we decided to restrict our analysis to the cerebral cortex of a single hemisphere. To this end, we used the same structural imaging data, pre-processing and probabilistic tracking pipeline as reported by Finger et al., but restricted subsequent processing to the 33 regions of interest (ROIs) of the left hemisphere. This data set included diffusion- and T1-weighted images acquired from 17 healthy subjects (7 female, age mean = 65.6y ± 10.9y) with a 3 Tesla Siemens Skyra MRI scanner (Siemens, Erlangen, Germany) and a 32-channel head coil. The 33 ROIs were registered individually for each subject based on the 'Desikan-Killiany' cortical atlas available in the Freesurfer toolbox (surfer.nmr.mgh.harvard.edu). This gave us the euclidean distances between each pair of ROIs. The incoming connections to each

Table 2. Model Parameters
region were normalized such that they summed up to 1. Since we were only interested in synchronization along indirect pathways, we needed some connections in our model to be strictly 0. Otherwise, it would be difficult to exclude potential synchronization along very weak direct connections. Hence, we chose to set all connections below a strength of 0.1 to zero. Afterwards, we re-normalized the input to each region such that they summed up to 1. The resulting SC matrix as well as the pair-wise distances are visualized in Fig 4A and 4B in the main paper.

Empirical Resting-State Functional Connectivity

Based on those SC and distance information we aimed to build a model of cortical activity able to reflect empirically observed synchronization behavior. Thus we needed empirical observations of cortical activity to evaluate our model. For this purpose, we acquired EEG data from the same 17 subjects as described above. This was done with 63 cephalic active surface electrodes arranged according to the 10/10 system (actiCAP R Brain Products GmbH, Gilching, Germany) for eight minutes of eyes-open resting-state. Again, data acquisition and pre-processing followed the same procedure as reported by Finger et al. [19]. EEG time-series from the surface electrodes were projected onto the centers of the ROIs via a linear constraint minimum variance spatial beam former [35]. The resulting source-space signals were band-pass filtered at 10 Hz and turned into analytic signals using the Hilbert transform. Subsequently, functional connectivity was evaluated as the coherence between all pairs of ROIs [36]. This resulted in the 33 x 33 functional connectivity matrix that can be observed in Fig 4C in the main paper and served as optimization target for our model.

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