Air-Core-Liquid-Ring (ACLR) Atomization: Influences of Gas Pressure and Atomizer Scale Up on Atomization Efficiency

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Received: 14 November 2018; Accepted: 22 February 2019; Published: 6 March 2019

Abstract: Air-core-liquid-ring (ACLR) atomizers present a specific type of internal mixing pneumatic atomizers, which can be used for efficient atomization of high viscous liquids. Generally, atomization efficiency is considered as a correlation between energy input and resulting droplet size. In pneumatic atomization, air-to-liquid ratio by mass (ALR) is commonly used as reference parameter of energy input. However, the pressure energy of the atomization gas is not considered in the calculation of ALR. In internal mixing ACLR atomizers, it can be assumed that this energy contributes to liquid disintegration by expansion of the gas core after exiting the atomizer. This leads to the hypothesis that droplet sizes decrease with increasing gas pressure at constant ALR. Therefore, the use of volumetric energy density ($E_V$) as a reference parameter of energy input was investigated at different gas pressures between 0.4 and 0.8 MPa. Furthermore, scale up-related influences on the atomization efficiency of ACLR atomization were investigated by use of an atomizer with enlarged exit orifice diameter. We can conclude that $E_V$ can be applied as a reference parameter of ACLR atomization processes with different gas pressures. However, within the range investigated no clear influence of gas pressure on atomization efficiency was found. Up-scaling of ACLR atomizers allows production of similar droplet sizes, but atomization efficiency decreases with increasing exit orifice diameter.

Keywords: ACLR; pneumatic atomization; high viscous feeds; efficiency; scale up

1. Introduction

Continuous atomization of high viscous concentrates is an important task in different industrial processes, like combustion, spray drying or coating [1,2]. In particular in spray drying, the capability to atomize high viscous feeds into small and uniform droplets is of great interest, as it allows an increase the initial dry matter content prior to actual drying [3,4].

For atomization of high viscous feeds, external-mixing pneumatic (ExMP) atomizers are commonly used [2]. In pneumatic atomizers the kinetic energy of a compressed gas stream, mostly air, is used as source of atomization energy [5–7]. In ExMP atomizers, gas and liquid are mixed under atmospheric conditions, outside the atomizer. As both media are supplied in separated streams, ExMP atomizers provide easy and independent control of gas and liquid flow rates [8]. Consequently, ExMP atomizers are widely used in lab and pilot scale spray drying. However, especially in spray-drying processes they do not meet the requirements of industrial applications on atomization efficiency, due to very high gas consumption rates [3,4]. In this case atomization efficiency is defined as a correlation between energy input and resulting spray droplet size. In pneumatic atomization, the energy input is typically correlated with the specific gas consumption under applied processing conditions, which is
expressed as the air-to-liquid ratio (ALR) by mass (Equation (1)). Here, \( M_X \) are the mass flows, \( Q_X \) are volume flows and \( \rho_X \) are densities of gas or liquid (L).

\[
ALR = \frac{M_{\text{gas}}}{M_L} = \frac{Q_{\text{gas}} \rho_{\text{gas}}}{Q_L \rho_L}
\]  

(1)

For ExMP atomizers, ALR in a range from 1 to 10 is usually applied [1,3]. The specific atomization efficiency of an atomizer can be estimated by use of a suitable correlation function [1]. Over the years, a variety of different correlation models has been developed for all kinds of pneumatic atomizers [9]. Generic models rely on the weighting of basic parameters, like nozzle exit diameter, processing conditions and liquid properties [10–12]. More recent correlation functions also contain so-called nozzle constants, which can only be determined experimentally [13,14]. Consequently, related correlations can be considered as “not complete”, as additional work is required for every change of the atomizer geometry. However, nozzle constants ensure the validity of correlations for a wide range of properties and nozzle geometries. A very popular correlation function for ExMP atomizers was proposed by Harari and Sher [13] (Equation (2)).

\[
SMD = B \cdot (\text{We} \cdot ALR)^{-0.4}
\]  

(2)

Here \( B \) is a characteristic nozzle constant and \( \text{We} \) is the Weber number of the gas. The Weber number represents the ratio between the air inertia forces and surface tension forces and is defined by Equation (3). In this calculation, \( p_{\text{Gas}} \) is the gas density, \( W_i^2 \) is the relative velocity, \( d_0 \) is the exit orifice diameter of the nozzle and \( \sigma_L \) is the surface tension of the liquid.

\[
\text{We} = \frac{p_{\text{gas}} W_i^2 d_0}{\sigma_L}
\]  

(3)

In contrast to ExMP atomizers, internal mixing pneumatic (IMP) atomizers are known for high atomization efficiency, as they allow for atomization of high viscous feeds at reduced ALR [15–17]. Consequently, they are considered to have high potential for use in industrial spray drying [18,19]. In IMP atomizers, atomization gas and liquid stream are brought into contact in a mixing zone inside the atomizer and leave the atomizer through the exit orifice in a two-phase flow pattern. For this purpose, several gas injection geometries have been proposed (e.g., [15,17,20–23]). Although no targeted mixing of gas and liquid is performed in most cases [24], it is well known for IMP atomizers that internal flow conditions determine the means of liquid disintegration and the spray development in time and space [22]. Especially an annular flow pattern inside the exit orifice is considered to be beneficial for stable atomization results [25,26].

The air-core-liquid-ring (ACLR) atomizer was designed especially to induce an annular flow pattern in the exit orifice by injection of a continuous core of compressed gas in the middle of the liquid stream [27]. In contrast to ExMP atomizers, atomization gas gets into contact with the liquid in a compressed state. By leaving the atomizer via the exit orifice, it expands and reaches atmospheric pressure. Inside the exit orifice, gas properties depend on the specific pressure loss of the atomizer, which in turn depends on the atomizer geometry, the ALR and the viscosity of the liquid. This is mainly based on the fact that the gas and the liquid stream have to share the cross-sectional area of the exit orifice in annular flow [28]. Consequently, ALR can be increased by a decrease in liquid flow rate (\( Q_L \)) at constant gas pressure (\( p_{\text{gas}} \)), or by an increase in \( p_{\text{gas}} \) at constant liquid flow rate (\( Q_L \)) [29]. In the latter case, the mass flow of gas (\( M_{\text{gas}} \)) increases at constant volume flow (\( Q_{\text{gas}} \)) by an increased gas density (\( \rho_{\text{gas}} \)). In practical use, this case of ALR adjustment is very important, as subsequent processes often require the adjustment of certain droplet size distributions at constant liquid flow rate and constant atomizer geometry. An increase in \( p_{\text{gas}} \) also increases the internal energy of the atomization gas stream by means of a higher expansion potential. This additional energy might contribute to disintegration of the liquid flow into droplets. This consideration is specific to
IMP atomizers. In ExMP atomizers, the influence of $p_{gas}$ can be neglected, as full gas expansion to atmospheric conditions can be assumed before the first contact with the gas stream. Consequently, velocities of gas and liquid are determined by atomizer geometry and ALR in ExMP atomizers. This is not the case for IMP atomizers, as the share of cross-sectional occupation of the exit orifice for liquid and gas is dependent on the pressure-dependent gas density. Accordingly, the aforementioned correlation function by Harari and Sher [13] is not applicable for IMP atomizers without knowledge of the gas pressure distribution inside the exit orifice, as the velocities of gas and liquid are unknown and change over the length of the exit orifice. As the expansion energy of the gas stream is not directly considered in the calculation of ALR (see Equation (1)) it is possible to apply equal ALR at different gas pressures with different gas expansion potentials in IMP atomizers. Based on the assumption that the energy of gas expansion contributes to the atomization energy, we hypothesize that spray droplet sizes decrease with increasing $p_{gas}$ at constant ALR. In this case, ALR is not sufficiently defined to represent the energy input of IMP atomizers. However, it is important to keep in mind that an increase in gas pressure is always connected to a change in gas density. Consequently, an increase in gas pressure always leads to an increase in ALR, when liquid flow rate and atomizer geometry are kept constant in IMP atomization. Therefore, it is not possible to investigate the sole influence of gas pressure on the atomization efficiency of an IMP atomizer with given geometry at constant ALR. Therefore, $p_{gas}$ and $Q_L$ have to be increased simultaneously in a specific ratio to keep the ALR constant. Hence, the influence of $p_{gas}$ can only be estimated in a comparative way, using an alternative reference parameter for quantification of the energy input.

As an alternative parameter of energy input, the volume-specific energy density ($E_V$) can be used. This parameter was originally proposed to evaluate the energy input of emulsification processes in different dispersing systems. According to Equation (4) it is defined as power input ($P$) per volume flow of liquid ($Q_L$) [30].

$$E_V = \frac{P}{Q_L}$$

(4)

As an alternative correlation function to estimate the atomization efficiency, a characteristic droplet size $x_{char}$ can be described as a function of $E_V$ by a power-law equation with two fit parameters $b$ and $c$ (Equation (5)) [30].

$$x_{char} = c \cdot E_V^{-b}$$

(5)

The characteristic droplet size is mostly expressed as Sauter mean diameter (SMD). However, other characteristic values of the generated droplet size distribution can also be used, depending on the objective of the investigation. To characterize the breakup of a bulk phase at a specific energy input, the characteristic droplet size $x_{90.3}$ can be used, representing the largest droplets of a distribution [31].

Stähle et al. introduced the volume specific energy density ($E_V$) in the field of atomization, in order to compare the atomization efficiency of atomizers that use different energy sources [32]. According to the authors, this value is very useful for atomizer selection and scale up of atomization processes [32]. In contrast to the ALR, which is only applicable in pneumatic atomization, $E_V$ is defined as a discrete value (power per volume flow of processed liquid). This allows direct evaluation of energy consumptions in different atomization processes, independent of the used atomizer.

For internal mixing atomizers $E_V$ is calculated according to Equation (6), where $R$ is the gas constant for dry air, $T$ is the temperature, and $p_0$ is the ambient pressure.

$$E_V = p_{gas} + \frac{M_G \cdot R \cdot T \cdot \ln\left(\frac{p_{gas} + p_0}{p_0}\right)}{Q_L}$$

(6)

In contrast to the ALR, the expansion potential of the gas stream is considered in $E_V$ in form of the energy that is needed for isothermal gas compression [33]. In order to review the correlation of
ALR and $E_V$, Equations (1) and (4) can be combined. Therefore, the liquid volume flow rate ($Q_L$) is expressed as a quotient of liquid mass flow rate ($M_L$) and liquid density ($\rho_L$). This leads to Equation (7).

$$E_V = p_{\text{gas}} + \frac{M_G}{M_L} \cdot \rho_L \cdot R \cdot T \cdot \ln \left( \frac{p_{\text{gas}} + p_0}{p_0} \right)$$  \hspace{1cm} (7)

Consequently, the mass flow ratio of gas and liquid can be replaced by the ALR (Equation (8)):

$$E_V = p_{\text{gas}} + \text{ALR} \cdot \rho_L \cdot R \cdot T \cdot \ln \left( 1 + \frac{p_{\text{gas}}}{p_0} \right)$$  \hspace{1cm} (8)

Based on this correlation, it can be seen that $E_V$ increases with increasing gas pressure ($p_{\text{gas}}$) at constant ALR, temperature ($T$) and liquid density ($\rho_L$). Therefore, as described above the increase of internal energy of the gas stream at increased $p_{\text{gas}}$ is considered in the calculation of $E_V$.

Nevertheless, it is not yet clear to which extent this energy is transformed into atomization energy in ACLR atomizers. Investigations of Stähle et al. were performed at constant $p_{\text{gas}}$ of 0.4 MPa. To the best of our knowledge, no further work on the applicability of $E_V$ as a reference parameter of ACLR atomization efficiency has been published yet. However, it can be assumed that differences in spray droplet sizes, caused by variation of $p_{\text{gas}}$ at constant ALR, are reflected by the value of $E_V$. This assumption can apply, if the portion of energy that is converted into atomization energy is larger than the energy of thermal losses. Thermal losses are not reflected by the described calculation of $E_V$, as the energy of the gas stream is calculated on the basis of isothermal gas compression.

In order to investigate the hypothesis of decreasing droplet sizes with increasing gas pressure, we performed atomization experiments with a pilot scale ACLR atomizer at constant $p_{\text{gas}}$ and different $Q_L$, as well as at constant $Q_L$ and different $p_{\text{gas}}$. In these experiments, model solutions of different dry matter concentrations and hence different viscosities were used. From all processing conditions, ALR and $E_V$ were calculated. Thereafter, resulting droplet sizes were correlated to ALR, as well as to $E_V$, in order to identify the applicability of both parameters as the reference of energy input in ACLR atomization. Additionally, the proposed advantages of using $E_V$ in the evaluation of scale up-related differences in atomization efficiency were investigated. As described above, energy efficiency of an ACLR atomizer depends on the pressure loss over the exit orifice, which in turn is dependent on geometry and processing conditions. Therefore, an up-scaled ACLR atomizer with larger exit orifice diameter was designed and used in atomization experiments with increased $Q_L$. Finally, the atomization efficiencies of both atomizers were compared.

2. Materials and Methods

Aqueous solutions of maltodextrin (C*Dry MD01958, Overlack GmbH, Mönchengladbach, Germany) were used as model system in the performed investigations. Dry matter contents $c_{\text{MD}}$ of 40, 45% and 47% were adjusted under consideration of the moisture content ($X = 4.5\%$, dry basis) of the used maltodextrin powder. The maltodextrin solutions were characterized by viscosity ($\mu$), surface tension ($\sigma$), density ($\rho$) and refractive index (RI). All measurements were executed in triplicate at a temperature of 25 °C. Mean values and relative uncertainty values $u_v(\gamma)$ of these parameters are summarized in dependency of the dry matter content $c_{\text{MD}}$ in Table 1. Viscosity measurements were performed using a rotary rheometer (MCR 101/301, Anton Paar GmbH, Graz, Austria) with coaxial cylinder geometry (CC27). Within the measurement range of shear-rates between 1 and 1000 s$^{-1}$ all solutions showed Newtonian behavior. Surface tensions were measured with a Wilhelmy plate system (DCAT 21, DataPhysics Instruments GmbH, Filderstadt, Germany). For measurement of refractive indices (RI) a refractometer (Carl-Zeiss, Oberkochen, Germany) was used. Densities were measured by use of a pycnometer.
Table 1. Physical properties viscosity ($\mu$), surface tension ($\sigma$), density ($\rho$) and refractive index (RI), as well as corresponding uncertainty values $u_r(\gamma)$ of the model solutions at 25 °C. Dry matter contents $c_{MD}$ are given as dry base values (d.b.).

| $c_{MD}$ | $\mu$ (Pa·s) | $u_r(\mu)$ (%) | $\sigma$ (N·m$^{-1}$) | $u_r(\sigma)$ (%) | $\rho$ (kg·m$^{-3}$) | $u_r(\rho)$ (%) | RI (-) | $u_r(RI)$ (%) |
|---------|-------------|----------------|----------------------|------------------|----------------------|-----------------|--------|--------------|
| 40      | 0.12        | 9.79           | 0.073                | 0.02             | 1175                 | 0.02            | 1.40   | 0.07         |
| 45      | 0.30        | 5.83           | 0.074                | 0.06             | 1206                 | 0.03            | 1.41   | 0.06         |
| 47      | 0.39        | 3.79           | 0.074                | 0.05             | 1218                 | 0.06            | 1.42   | 0.04         |

2.1. Air-Core-Liquid-Ring Atomizers

In this study, two different self-manufactured ACLR atomizers prototypes were used. The geometries of both atomizers are shown in Figure 1. Atomizer 1 is a prototype, designed for atomization in pilot scale spray-drying processes. This geometry was already used in several former studies [27,29,32,34–36]. Atomizer 2 is a scale up version of atomizer 1. In comparison to atomizer 1, mainly diameter and length of the exit orifice were increased. The most important dimensions of both atomizers are given in Table 2.

Figure 1. Schemes of used air-core-liquid-ring atomizers 1 and 2. Dimensions of gas capillary diameter (A) and length of the mixing zone (B), as well as length (C) and diameter (D) of the exit orifice are given in Table 2.

Table 2. Dimensions of the used Air-core-liquid-ring (ACLR) atomizers: gas capillary diameter (A), length of the mixing zone (B), length of the exit orifice (C) and diameter of the exit orifice (D).

| Dimension | Atomizer 1 | Atomizer 2 |
|-----------|------------|------------|
| A         | 1.5 mm     | 1.5 mm     |
| B         | 2.4 mm     | 2.4 mm     |
| C         | 1.5 mm     | 2.5 mm     |
| D         | 1.5 mm     | 3 mm       |

2.2. Spray Test Rig

For spray experiments a modular test rig was used. The atomization gas was compressed air supplied by a compressor (Renner RSF-Top 7.5, Renner GmbH, Güglingen, Germany) with a pressure vessel volume of 90 L. Relative pressures of 0.4, 0.6, and 0.8 MPa were adjusted with a pressure regulator. Corresponding gas volume flows ($Q_{gas}$) were measured by a gas flow meter (ifm SD6000, ifm electronic, Essen, Germany). Measurement values of $Q_{gas}$ are reported at normal conditions (273.15 K, 0.1013 MPa$\text{abs}$) in normal cubic meters per hour (Nm$^3$h$^{-1}$). In experiments with atomizer 1, an eccentric screw pump (NM011BY, Erich Netzsch GmbH & Co. Holding KG, Waldkraiburg, Germany) was used for the liquid supply. The flow rate was adjusted between 21 L·h$^{-1}$ and 40 L·h$^{-1}$ by an electronic frequency converter. In experiments with atomizer 2, a pressure vessel was used to supply flow rates of 100 L·h$^{-1}$, 120 L·h$^{-1}$, and 140 L·h$^{-1}$, adjusted by a manually operated needle.
In both cases, the liquid flow rate was measured by a flow meter (VSI 044/16, VSE GmbH, Neuenrade, Germany). The liquid pressure was not monitored. All trials were performed at room temperature (25 °C).

2.3. Droplet Size Measurement

For spray droplet size measurements a laser diffraction spectroscope (Spraytec, Malvern Instruments, Malvern, UK) was mounted to the test rig. The spectroscope was equipped with a 750 mm focal lens, offering a droplet size measurement range of 2–2000 µm. The laser beam crossed the full cone of the spray angle at the atomizers centerline in a vertical distance of 250 mm. The spray was collected in a vessel below the measurement zone, which was connected to an exhaust fan. Measurements were conducted at a frequency of 250 Hz over a time of 25 s, leading to the recording of 6250 droplet size distributions per measurement. Based on this data, an average droplet size distribution was calculated under consideration of the refractive indices of the used model solutions. For beam steering correction, the method described in one of our former studies was used [36]. In this study, the commonly used Sauter mean diameter (SMD) was chosen as standard reference parameter. Additionally, the characteristic value \( x_{90.3} \) was regarded to review the correlation between energy input and resulting droplet sizes. This parameter represents the largest droplets of a distribution and is known to be very sensitive to changes in processing conditions of ACLR atomization. Results of \( x_{90.3} \) are only partly shown, in order to ensure readability of the paper.

2.4. Calculations and Statistical Analysis

OriginPro software, version 2017G (OriginLab Corporation, Northampton, MA, USA) was used for calculation of fitting curves and statistical analysis. For fitting curves, the Levenberg Marquardt iteration algorithm was used. As goodness of fit parameters, the corrected determination coefficient \( R^2 \) was used. The effects of ALR and \( E_V \) on resulting spray droplet sizes were evaluated in dependency of the liquid viscosity by a 1-way analysis of variance (ANOVA). Scheffé’s test was used for comparison of means. In the performed tests, probability of \( p < 0.05 \) was used for the identification of significant differences.

3. Results and Discussion

3.1. Correlation between Volumetric Energy Density (\( E_V \)) and Air-to-Liquid Ratio (ALR) of Applied Processing Conditions

In this section, the correlation between the standard reference parameter of energy input ALR and the potential alternative reference parameter \( E_V \) for ACLR atomization is reviewed. For calculation of \( E_V \) in dependency of ALR (Equation (8)), a temperature (\( T \)) of 25 °C and an atmospheric pressure (\( p_0 \)) of 0.1013 MPa was assumed. Liquid densities (\( \rho_L \)) of model solutions only vary in a small range (see Table 1). Therefore, the arithmetic mean value \( \rho_{L, av} = 1200 \text{ kg m}^{-3} \) was used for linear prediction of the correlation between ALR and \( E_V \) at constant gas pressure, based on Equation (8). These correlations are shown in Figure 2 by dashed lines. Moreover, processing conditions of performed experiments with atomizer 1 are indicated in Figure 2. For these calculations, \( \rho_L \) of each model solution was used. Therefore, slight differences from predictions on base of \( \rho_{L, av} \) occur.
The range of ALR is 0.05 to 0.35 in this case. These correlations are well known for internal mixing atomizers [1,36]. For each model solution, \( Q_{gas} \) increases with decreasing \( Q_L \) at constant \( p_{gas} \). Moreover, \( Q_{gas} \) increases at constant \( Q_L \) with increasing \( p_{gas} \), due to higher gas compression rates. Hence, ALR increases with increasing \( p_{gas} \) at constant \( Q_L \) and decreases with increasing \( Q_L \) at constant \( p_{gas} \). At a viscosity of 0.39 Pa·s, ALR ranges from 0.07 (\( p_{gas} = 0.4 \) MPa; \( Q_L = 30 \) L·h\(^{-1}\)) to 0.31 (\( p_{gas} = 0.8 \) MPa; \( Q_L = 21 \) L·h\(^{-1}\)). A decrease of \( \mu \) to 0.30 Pa·s leads to ALR between 0.08 and 0.31, when the aforementioned process conditions are applied. When the solution with \( \mu = 0.12 \) Pa·s is used, a flow rate of \( Q_L = 40 \) L·h\(^{-1}\) at \( p_{gas} = 0.4 \) MPa was applicable to the used pump. This leads to a decrease of the smallest ALR. The range of ALR is 0.05 to 0.35 in this case. These correlations are well known for internal mixing atomizers [1,36]. For \( E_V \), the same tendencies are present, as the same values of \( Q_L \) and \( Q_{gas} \) are used for calculation. The ranges of \( E_V \) are as follows: \( \mu = 0.39 \) Pa·s, \( E_V = 11.3–70.5 \); \( \mu = 0.30 \) Pa·s, \( E_V = 13.0–71.5 \); \( \mu = 0.12 \) Pa·s, \( E_V = 9.0–77.9 \).

**Figure 2.** Linear prediction of the correlation between \( E_V \) and air-to-liquid ratio (ALR) at constant \( p_{gas} \) and \( \rho_L \) (dashed lines). Additionally, applied processing conditions are indicated: \( p_{gas} = 0.4 \) MPa (squares), 0.6 MPa (triangles), and 0.8 MPa (circles); \( \mu = 0.12 \) Pa·s (empty symbols), 0.30 Pa·s (half-filled symbols), and 0.39 Pa·s (filled symbols).

From this illustration, it can be seen that process conditions with identical ALR and different \( E_V \), as well as with identical \( E_V \) and different ALR are applicable. Moreover, differences between both parameters increase with increasing \( p_{gas} \), as the slope of linear correlation increases. However, as the gas volume flow rate (\( Q_{gas} \)) cannot be set independently of \( p_{gas} \) and \( Q_L \), the range of overlapping ALRs at different \( p_{gas} \) is very small in applied processing conditions. For a detailed review of processing conditions, resulting gas volume flow rates (\( Q_{gas} \)), as well as calculated values of ALR and \( E_V \) are shown in dependency of \( p_{gas} \) and \( Q_L \) for the used model solutions (\( c_{MD} \), \( \mu \)) in Table 3.
Table 3. Gas pressures (\(p_{\text{gas}}\)), liquid volume flow rates (\(Q_L\)), resulting gas volume flow rates (\(Q_{\text{gas}}\)), air/liquid mass ratios (ALR) and volumetric energy densities (\(E_V\)) for atomizer 1 together with the uncertainty values \(u_r(\gamma)\) of measured/calculated parameters \(Q_{\text{gas}}\), ALR, and \(E_V\). Measurement values of \(Q_{\text{gas}}\) are reported at normal conditions (273.15 K, 0.1013 MPa abs) in normal cubic meters per hour (Nm\(^3\)·h\(^{-1}\)).

| Atomizer 1 |  |  |  |  |  |  |  |
|------------|---|---|---|---|---|---|---|
| c\(_{MD}\) | \(\mu\) | \(p_{\text{gas}}\) | \(Q_L\) | \(Q_{\text{gas}}\) | \(u_r(Q_{\text{gas}})\) | ALR | \(u_r(\text{ALR})\) | \(E_V\) | \(u_r(E_V)\) |
| (%) d.b. | (Pa·s) | (Mpa) | (L·h\(^{-1}\)) | (Nm\(^3\)·h\(^{-1}\)) | (%) | (-) | (%) | (kJ·L\(^{-1}\)) | (%) |
| 0.4 | 21 | 2.9 | 2.01 | 0.15 | 3.57 | 24.1 | 3.72 |
| | 30 | 2.5 | 2.34 | 0.09 | 4.08 | 14.8 | 4.15 |
| | 40 | 1.9 | 2.99 | 0.05 | 5.40 | 9.0 | 5.30 |
| 0.6 | 21 | 4.7 | 2.13 | 0.24 | 2.22 | 47.9 | 2.45 |
| | 30 | 4.1 | 1.40 | 0.15 | 2.55 | 30.0 | 2.72 |
| | 40 | 3.5 | 2.86 | 0.10 | 3.37 | 19.4 | 3.44 |
| 0.8 | 21 | 6.8 | 2.55 | 0.35 | 1.52 | 77.9 | 1.81 |
| | 30 | 6.0 | 0.96 | 0.22 | 1.68 | 49.3 | 1.93 |
| | 40 | 5.1 | 1.96 | 0.14 | 1.99 | 31.4 | 2.18 |
| 0.4 | 21 | 2.7 | 0.00 | 0.14 | 3.72 | 22.9 | 3.86 |
| | 25 | 2.1 | 2.71 | 0.09 | 4.69 | 15.6 | 4.73 |
| | 30 | 2.2 | 2.66 | 0.08 | 4.62 | 13.0 | 4.64 |
| 0.6 | 21 | 4.6 | 1.26 | 0.23 | 2.21 | 47.0 | 2.44 |
| | 30 | 3.7 | 1.55 | 0.13 | 2.69 | 26.9 | 2.84 |
| | 40 | 3.0 | 1.90 | 0.08 | 3.31 | 16.8 | 3.37 |
| 0.8 | 21 | 6.2 | 2.48 | 0.31 | 1.65 | 71.5 | 1.92 |
| | 30 | 5.5 | 2.11 | 0.19 | 1.85 | 44.0 | 2.07 |
| | 40 | 4.7 | 1.24 | 0.13 | 2.15 | 29.1 | 2.31 |
| 0.4 | 21 | 2.4 | 2.37 | 0.12 | 4.13 | 20.6 | 4.23 |
| | 23 | 1.9 | 2.99 | 0.09 | 5.18 | 15.2 | 5.19 |
| | 30 | 1.9 | 3.09 | 0.07 | 5.37 | 11.3 | 5.32 |
| 0.6 | 21 | 4.2 | 1.36 | 0.21 | 2.38 | 43.6 | 2.59 |
| | 30 | 3.4 | 0.00 | 0.12 | 3.01 | 24.2 | 3.13 |
| | 40 | 3.1 | 3.23 | 0.08 | 3.87 | 17.6 | 3.88 |
| 0.8 | 21 | 6.1 | 0.95 | 0.31 | 1.67 | 70.5 | 1.93 |
| | 30 | 5.0 | 0.00 | 0.17 | 2.62 | 40.0 | 2.75 |
| | 40 | 4.9 | 2.04 | 0.13 | 5.04 | 30.3 | 5.04 |

3.2. Droplet Sizes in Dependency of \(E_V\) and ALR

In Figure 3, SMD is shown in the dependency of ALR (left) and \(E_V\) (right). Results are presented according to the used liquid viscosity (\(\mu\)): 0.39 Pa·s (red, top line), 0.30 Pa·s (blue, middle line), and 0.12 Pa·s (black, bottom line). The applied \(p_{\text{gas}}\) is represented by different symbols: 0.4 MPa (empty), 0.6 MPa (half-filled), and 0.8 MPa (filled). Atomization efficiency is expressed by the following fit functions (Equations (9) and (10)), which are adapted from Equation (5):

\[
\text{SMD} = c \times ALR^{-b} \tag{9}
\]

\[
\text{SMD} = c \times E_V^{-b} \tag{10}
\]
Figure 3. Characteristic droplet size Sauter mean diameter (SMD) in dependence of ALR (left) and volumetric energy density ($E_V$) (right) for viscosities ($\mu$) of 0.12 Pa·s (black), 0.30 Pa·s (blue), and 0.39 Pa·s (red) at $p_{gas}$ of 0.4 MPa (empty triangles), 0.6 MPa (half-filled circles), and 0.8 MPa (filled squares). Additionally, fits of results at the same $\mu$ and same $p_{gas}$ are depicted by dashed lines. Total fits of all results at the same $\mu$ and different $p_{gas}$ are depicted by solid lines.

As described above, $R^2$ is used to review the goodness of fit. Fits of results at the same viscosity ($\mu$) and same $p_{gas}$ are depicted by dashed lines. Within each fit, the level of gas compression is constant. Therefore, no differences in values of $R^2$ are expected between applications of ALR and $E_V$ as a reference parameter, when results of same gas pressure and same viscosity are compared. Total fits of all results at same $\mu$ but different $p_{gas}$ are depicted by solid lines. In this case, higher $R^2$ values are expected when $E_V$ is used as a reference parameter, as the effect of gas compression has to be represented by the applied fit. An overview of fit parameters and determination coefficients $R^2$ is given in Table 4.
The overall dependencies between droplet sizes and both reference parameters are very similar. For all viscosities, SMD decreases with increasing energy input (ALR or $E_V$), in a non-linear way. At low energy inputs, SMD decreases strongly with increasing energy input. With increasing energy input, this decrease is lowered. An increase in viscosity at constant energy input results in a significant increase of SMD. These tendencies are described well in literature for the correlation between ALR and SMD [20,27,37], and were also found for the correlation between energy input and the characteristic droplet size $x_{90,3}$ (data not shown). At a viscosity of 0.39 Pa·s, SMD decreases from 106 µm (ALR = 0.07; $E_V = 11.3$ kJ·L$^{-1}$) to 41 µm (ALR = 0.31; $E_V = 70.5$ kJ·L$^{-1}$). Accordingly, at a viscosity of 0.30 Pa·s SMD decreases from 80 µm (ALR = 0.08; $E_V = 13.0$ kJ·L$^{-1}$) to 29 µm (ALR = 0.31; $E_V = 71.5$ kJ·L$^{-1}$). In case of a viscosity of 0.12 Pa·s SMD decreases from 81 µm (ALR = 0.05; $E_V = 9.0$ kJ·L$^{-1}$) to 27 µm (ALR = 0.35; $E_V = 77.9$ kJ·L$^{-1}$).

As expected, no significant differences in determination coefficients $R^2$ of correlations between SMD and the both reference parameters ALR and $E_V$ are present at constant $p_{gas}$ (0.4, 0.6, or 0.8 MPa) and constant viscosity (0.12, 0.30, or 0.39 Pa·s) (see Table 4). The fit functions used represent well the correlations between SMD and the reference parameters over the whole range of investigated energy inputs. However, at a viscosity above 0.30 Pa·s $R^2$ decreases from 0.99 at $p_{gas}$ of 0.6 MPa to 0.80 at $p_{gas}$ of 0.4 MPa. In this case, the deviation between the applied fit function and measurement results increases with decreasing $p_{gas}$. It can be concluded that an increase in viscosity leads to a higher demand of energy input, in order to create an annular flow pattern in the exit orifice of an ACLR atomizer. Hence, $p_{gas}$ of 0.4 MPa might not be sufficient for the setup of a stable annular flow, at viscosities above 0.30 Pa·s. However, further research is needed to evaluate this assumption.

When all results of the same viscosity are subjected to a total fit, larger differences in $R^2$ are present, when results for ALR and $E_V$ are compared. In the case of ALR, $R^2$ is identical for viscosities of 0.12 and 0.30 Pa·s (0.96) and decreases to 0.90 at a viscosity of 0.39 Pa·s. In case of $E_V$, $R^2$ increases with increasing viscosity from 0.80 (0.12 Pa·s) to 0.92 (0.30 Pa·s) and 0.97 (0.39 Pa·s). Consequently, in case of $E_V$ the deviation between measurement points and fit equations decreases with increasing viscosity, while the deviation increases in the case of ALR with an increase in viscosity above 0.30 Pa·s.

In conclusion, the effect of gas expansion might also increase with increasing viscosity. However, within the investigated parameter range the influence of gas pressure on atomization efficiency cannot be determined to a full extent. In contrast to ALR, the expansion energy of the gas stream is considered in a calculation of $E_V$. However, in which way and to which extent the energy contributes to droplet

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**Table 4.** Fit parameters $b$ and $c$, as well as determination coefficient $R^2$ for correlations between Sauter mean diameter (SMD) and ALR, as well as between SMD and $E_V$.

| Atomizer 1 | $c_{MD}$ (%) d.b. | $\mu$ (Pa·s) | $p_{gas}$ (Mpa) | SMD = $c \times ALR^{-b}$ | SMD = $c \times E_V^{-b}$ |
|------------|-------------------|--------------|-----------------|-----------------------------|-----------------------------|
|            |                   |              |                 | $b$ | $c$ | $R^2$ | $b$ | $c$ | $R^2$ |
| 40         | 0.12              | 0.4          | 0.685           | 11.3 | 0.96 | 0.706 | 396 | 0.96 |
|            |                   | 0.6          | 0.776           | 11.2 | 0.90 | 0.793 | 720 | 0.90 |
|            |                   | 0.8          | 0.506           | 15.9 | 0.99 | 0.514 | 254 | 0.99 |
|            | **total**         |              |                 | 0.634 | 12.9 | 0.96 | 0.321 | 163 | 0.80 |
| 45         | 0.30              | 0.4          | 0.597           | 16.2 | 0.80 | 0.611 | 361 | 0.80 |
|            |                   | 0.6          | 0.695           | 13.4 | 0.97 | 0.711 | 573 | 0.97 |
|            |                   | 0.8          | 0.805           | 12.4 | 0.93 | 0.821 | 1048 | 0.93 |
|            | **total**         |              |                 | 0.706 | 13.2 | 0.96 | 0.680 | 514 | 0.92 |
| 47         | 0.39              | 0.4          | 0.968           | 7.1  | 0.80 | 0.994 | 1107 | 0.80 |
|            |                   | 0.6          | 0.516           | 20.8 | 0.99 | 0.528 | 339 | 0.99 |
|            |                   | 0.8          | 0.378           | 26.2 | 0.99 | 0.384 | 210 | 0.99 |
|            | **total**         |              |                 | 0.413 | 24.7 | 0.90 | 0.531 | 265 | 0.97 |
disintegration is not considered. Possibly, the gas core mainly expands inside the exit orifice. In this case, expansion energy would mainly be transformed into kinetic energy. Consequently, an increased gas pressure would lead to an acceleration of the gas core and, therefore, to the transmission of higher amounts of energy to the liquid sheet. In turn, this would result in a viscosity dependent decrease of the liquid sheet thickness at the end of the atomizer. Consequently, the direct correlation of energy input and resulting droplet sizes might not be sufficient for determination of the atomization efficiency. Instead, the liquid sheet thickness has to be taken into account for assessment of atomization efficiency, although it is difficult to evaluate in practical use. For investigation of this assumption, further studies on process dependent flow conditions inside the exit orifice of ACLR atomizers have to be carried out, in order to gain information on the relative velocity of the gas and the liquid phase.

When this information is available, the correlation function for ExMP atomizers, like that proposed by Harari and Sher [13] might be applicable for ACLR atomization as well.

3.3. Influence of Atomizer Scale Up on Atomization Efficiency

In order to evaluate scale up-related differences in atomization efficiency, additional trials with a scaled up version of the ACLR atomizer described above were performed (atomizer 2) at \( \mu = 0.12 \text{ Pa·s} \). In atomizer 2, the diameter of the exit orifice (\( d_{\text{exit orifice}} \)) was increased to 3.0 mm (see Section 2.2). Accordingly, liquid volume flow rates (\( Q_L \)) were increased to 100, 120, and 140 L·h\(^{-1} \), while the same \( p_{\text{gas}} \) as in the trials with atomizer 1 were applied (0.4, 0.6, and 0.8 MPa). In Table 5, gas volume flow rates (\( Q_{\text{gas}} \)) and corresponding uncertainty values (\( u_r(Q_{\text{gas}}) \)) are shown in a dependency with \( Q_L \). From this data, ALR and \( E_V \) are calculated according to Equations (1) and (6). Results and corresponding uncertainty values are also presented in Table 5.

Table 5. Applied gas pressures (\( p_{\text{gas}} \)) and liquid volume flow rates (\( Q_L \)), as well as resulting gas volume flow rates (\( Q_{\text{gas}} \)), air/liquid mass ratios (ALR) and volumetric energy densities (\( E_V \)). Additionally, uncertainty values \( u_r(\gamma) \) of resulting parameters \( Q_{\text{gas}}, \text{ ALR, and } E_V \) are given. Measurement values of \( Q_{\text{gas}} \) are reported at normal conditions (273.15 K, 0.1013 MPa\(_{\text{abs}} \)) in normal cubic meters per hour (Nm\(^3\)·h\(^{-1} \)).

| Atomizer 2 | \( \gamma \) | \( p_{\text{gas}} \) | \( Q_L \) | \( u_r(Q_L) \) | \( Q_{\text{gas}} \) | \( u_r(Q_{\text{gas}}) \) | ALR | \( u_r(\text{ALR}) \) | \( E_V \) | \( u_r(E_V) \) |
|------------|------------|----------------|-------------|----------------|----------------|----------------|-------|----------------|----------------|----------------|
|            | \( \% \text{ d.b.} \) | (Pa·s) | (MPa) | (L·h\(^{-1} \)) | (\%) | (Nm\(^3\)·h\(^{-1} \)) | (\%) | (\%) | (kJ·L\(^{-1} \)) | (\%) |
| 40         | 0.12       | 0.4      | 100     | 0.21            | 13.2 | 2.00 | 0.14 | 1.85 | 24.1            | 2.01 |
| 120        | 0.82       | 12.0    | 7.36    | 0.11            | 8.05 | 18.6 | 8.36 |
| 141        | 0.48       | 11.7    | 3.08    | 0.09            | 2.78 | 15.5 | 2.71 |
| 0.6        | 100        | 0.63    | 19.6    | 0.88            | 0.22 | 0.92 | 43.3 | 0.92 |
| 120        | 0.05       | 18.9    | 0.03    | 0.17            | 0.07 | 34.9 | 0.07 |
| 140        | 0.68       | 18.4    | 1.13    | 0.14            | 1.76 | 29.2 | 1.89 |
| 0.8        | 100        | 0.66    | 24.0    | 0.24            | 0.26 | 0.87 | 58.4 | 0.91 |
| 120        | 0.46       | 23.5    | 0.65    | 0.22            | 0.90 | 48.2 | 0.91 |
| 140        | 0.22       | 22.9    | 0.25    | 0.18            | 0.47 | 40.8 | 0.46 |

In comparison to atomizer 1, absolute values of \( Q_{\text{gas}} \) are significantly higher, when atomizer 2 is used. However, when minimum and maximum values of both atomizers are compared, the factor of increase is in the range between 3.5 and 6.1 and, therefore, similar to the increase of \( Q_L \) (factor 3.5–4.8). Hence, a wide range of overlapping ALR and \( E_V \) is present in the results of both atomizers. In order to review the efficiency of both atomizers, values of SMD and \( x_{90,3} \) are shown in dependency of ALR and \( E_V \) in Figure 4a,b. Correlations between SMD and ALR, as well as between SMD and \( E_V \), are approximated according to Equations (9) and (10). Fit parameters and determination coefficients \( R^2 \)
are given in Table 6. Correlations between $x_{90,3}$ and ALR/EV are approximated by replacing SMD by $x_{90,3}$ in Equations (9) and (10). Fit parameters and determination coefficients $R^2$ are given in Table 6.

**Figure 4.** (a). SMD in dependency of ALR (left) and $E_V$ (right) at $\mu = 0.12$ Pa·s, using atomizers with exit orifice diameters of 1.5 mm (atomizer 1, $Q_L = 20–40$ L·h$^{-1}$; black filled symbols) and 3.0 mm (atomizer 2, $Q_L = 100–140$ L·h$^{-1}$; red empty symbols). Applied gas pressures are 0.4 MPa (squares), 0.6 MPa (triangles), and 0.8 MPa (circles). 

(b) Characteristic droplet size $x_{0,3}$ in dependency of ALR (left) and $E_V$ (right) at $\mu = 0.12$ Pa·s, using atomizers with exit orifice diameters of 1.5 mm (atomizer 1, $Q_L = 20–40$ L·h$^{-1}$; black filled symbols) and 3.0 mm (atomizer 2, $Q_L = 100–140$ L·h$^{-1}$; red empty symbols). Applied gas pressures are 0.4 MPa (squares), 0.6 MPa (triangles), and 0.8 MPa (circles).
Table 6. (a) Fit parameters $b$ and $c$, as well as determination coefficient $R^2$ for correlations between SMD and $E_V$. (b) Fit parameters $b$ and $c$, as well as determination coefficient $R^2$ for correlations between $x_{90,3}$ and $E_V$.

| SMD (% d.b.) | $\mu$ (Pa·s) | Atomizer | $d_{exit \ orifice}$ (mm) | SMD = $c \times ALR^{-b}$ | SMD = $c \times E_V^{-b}$ |
|--------------|--------------|----------|--------------------------|---------------------------|---------------------------|
| 40           | 0.12         | 1        | 1.5                      | 0.634                     | 12.9                      | 0.96                      | 0.531                     | 266                       | 0.97                      |
|              |              | 2        | 3.0                      | 0.544                     | 25.3                      | 0.78                      | 0.420                     | 293                       | 0.80                      |

(b)

| SMD (% d.b.) | $\mu$ (Pa·s) | Atomizer | $d_{exit \ orifice}$ (mm) | $x_{90,3} = c \times ALR^{-b}$ | $x_{90,3} = c \times E_V^{-b}$ |
|--------------|--------------|----------|--------------------------|--------------------------------|--------------------------------|
| 40           | 0.12         | 1        | 1.5                      | 0.387                         | 105.6                         | 0.94                      | 0.334                     | 680.7                     | 0.95                      |
|              |              | 2        | 3.0                      | 0.615                         | 91.2                          | 0.90                      | 0.485                     | 1502.6                    | 0.95                      |

For atomizer 2, SMD decreases from 97 $\mu$m to 56 $\mu$m with increasing energy input, while $x_{90,3}$ decreases from 402 $\mu$m to 222 $\mu$m. In the case of atomizer 1, $x_{90,3}$ decreases from 331 $\mu$m to 149 $\mu$m. In the correlation between SMD and ALR, $R^2$ decreases from 0.96 to 0.78 with increasing exit diameter of the atomizer. Identical values of $R^2$ are present for the correlation between $x_{90,3}$ and ALR. For the correlation between SMD and $E_V$, $R^2$ decreases from 0.97 to 0.80 with increasing exit orifice diameter. When $x_{90,3}$ is considered, $R^2$ of the applied fit is equal (0.95) for both atomizers. This is an interesting fact, as it can be concluded that the correlation between SMD and ALR seems to be independent of the used reference parameter. In turn, the correlation between $x_{90,3}$ and $E_V$ seems to be the best choice for scale-up considerations. This might be based on the fact that $x_{90,3}$ is a characteristic droplet size which is directly derived from the droplet size distribution and not a calculated mean diameter, like SMD. However, these considerations should be proved by further research.

Another interesting point is that similar values of SMD and $x_{90,3}$ can be produced by both atomizers in the investigated range of energy inputs (ALR and $E_V$). However, the necessary energy input is significantly higher for atomizer 2. Therefore, atomization efficiency decreases with increasing $d_{exit \ orifice}$. Based on the performed approximations (see Table 6), the factor of increase averages 2.2, in order to produce the same values of $x_{90,3}$. With increasing $E_V$ and, therefore, decreasing $x_{90,3}$ this factor decreases from 2.6 at $x_{90,3} = 330 \mu$m to 1.7 at $x_{90,3} = 210 \mu$m. Based on these numbers, the application of several smaller atomizers would be more energy efficient than the scale up of one atomizer by increasing $d_{exit \ orifice}$, when $Q_L$ has to be increased in practical use.

In Section 3.2 it was assumed that droplet size is mainly dependent on the thickness of the liquid sheet at the end of the exit orifice and relative velocity of gas and liquid phase at constant viscosity. Assuming similar relative velocities for both atomizers at the interface between gas core and liquid sheet, also liquid sheet thickness should be similar, when same droplet sizes are produced by the two atomizers.

In the case of similar liquid sheet thickness, an increase in dexit orifice leads to an overproportional increase of the cross-sectional area, occupied by the gas core inside the exit orifice. However, the kinetic energy of the gas, travelling in the middle of the gas core, cannot be used for liquid disintegration. Therefore, a higher relative volume flow of gas is needed at increased dexit orifice, leading to higher demand of energy input and hence higher $E_V$, when atomizer 2 is used. Still, no information on relative velocities or liquid sheet thicknesses is available at this point, which is why further investigations have to be carried out on this topic.

4. Conclusions

In the presented study, the influences of gas pressure ($p_{gas}$) and atomizer scale up on atomization efficiency of air-core-liquid-ring (ACLR) atomization were investigated. It was assumed that
pressure-dependent expansion energy of atomization gas contributes to liquid disintegration in ACLR atomization. For this reason it was hypothesized that droplet size decreases with increasing $p_{\text{gas}}$ at constant ALR. In this case ALR would not be sufficient for distinct determination of processing conditions. Therefore, the volumetric energy density ($E_V$) was considered as an alternative reference parameter.

For correlations between SMD/$x_{90,3}$ and the two investigated reference parameters of energy input, ALR and $E_V$, only slight differences in determination coefficients $R^2$ of applied fit functions were found. In total, an influence of gas pressure on the atomization efficiency could not be determined to full extend in the range investigated. Instead, the mechanism of liquid disintegration seems to be similar to that of ExMP atomizers. Hence, resulting droplet sizes mainly dependent on the liquid sheet thickness at the outlet of the atomizer, the relative velocity of gas and liquid phase, and the viscosity of the liquid. Consequently, the use of established correlation functions for ExMP atomizers should be considered for description of the atomization performance of ACLR atomizers in future studies.

Experiments with an up-scaled ACLR atomizer revealed that similar droplet sizes can be generated with increasing liquid flow rates, when the diameter of the exit orifice is increased. However, in this case also specific energy consumption increases, leading to a decreased atomization efficiency as the kinetic energy of the gas, travelling in the middle of the gas core, cannot be used for liquid disintegration. As a conclusion, the direct correlation of energy input and resulting droplet sizes might not be sufficient for determination of atomization efficiency of differently scaled ACLR atomizers. Instead, the liquid sheet thickness has to be taken into account for assessment of atomization efficiency and scale up, although it is difficult to evaluate in practical use.

Further research will be addressed to the flow conditions inside the exit orifice of ACLR atomizers, in order to gain a better understanding of the mechanism of droplet formation. In particular, the dependency of the liquid sheet thickness and the relative velocity of gas and liquid phase has to be investigated for further development of ACLR atomization.

**Author Contributions:** Conceptualization, M.O.W. and V.G.; Methodology, M.O.W.; Validation, M.O.W., H.P.K. and V.G.; Investigation, M.O.W.; Writing-Original Draft Preparation, M.O.W.; Writing-Review and Editing, H.P.K. and V.G.; Visualization, M.O.W.

**Funding:** This research received no external funding.

**Acknowledgments:** The authors express their thanks to Andreas Reimer, Philip Junginger and Andrea Butterbrodt for experimental support.

**Conflicts of Interest:** The authors declare no conflict of interest.

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