ATLAS Sensitivity to Leptoquarks, $W_R$ and Heavy Majorana Neutrinos in Final States with High-$p_T$ Dileptons and Jets with Early LHC Data at 14 TeV proton-proton collisions

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Dilepton-jet final states are used to study physical phenomena not predicted by the standard model. The ATLAS discovery potential for leptoquarks and Majorana Neutrinos is presented using a full simulation of the ATLAS detector at the Large Hadron Collider. The study is motivated by the role of the leptoquark in the Grand Unified theories and the see-saw mechanism that could explain the masses of the observed neutrinos. The analysis algorithms are presented, background sources are discussed and estimates of sensitivity and the discovery potential for these processes are reported.

1. Introduction

The Large Hadron Collider (LHC) will soon open up a new energy scale that will directly probe for physical phenomena outside the framework of the Standard Model (SM). Many SM extensions inspired by Grand Unified theories introduce new, very heavy particles such as leptoquarks. Extending the SM to a larger gauge group that includes, e.g. Left-Right Symmetry (LRS) [1], could also explain neutrino masses via the see-saw mechanism. The LRS-based Left-Right Symmetric Model (LRSM) [10] used as a guide for presented studies, extends the electroweak gauge group of the SM from $\text{SU}(2)_L \times \text{U}(1)_Y$ to $\text{SU}(2)_L \times \text{SU}(2)_R \times \text{U}(1)_B$, and thereby introduces $Z'$ and right-handed $W$ bosons. If the LRS breaking in nature is such that all neutrinos become Majoranas, the LRSM predicts the see-saw mechanism [2] that elegantly explains the masses of the three light neutrinos.

2. Search for scalar leptoquarks

Leptoquarks (LQ) are hypothetical bosons carrying both quark and lepton quantum numbers, as well as fractional electric charge [3,4]. Leptoquarks could, in principle, decay into any combination of any quark and any lepton. Experiments limit their branching fractions to a lepton and a quark from the same SM generation [5]. Leptoquarks can either be produced in pairs by the strong interaction or in association with a lepton via the leptoquark-quark-lepton coupling. Figure 1 shows Feynman diagrams for the pair production of leptoquarks at the LHC.

This contribution describes the search strategy for leptoquarks decaying to either an electron and a quark or a muon and a quark leading to final states with two leptons and at least two jets. The branching fraction of a leptoquark to a charged lepton and a quark is denoted as $^1$.

MC-simulated signal events have been studied using Monte Carlo (MC) samples for first generation (1st gen.) and second generation (2nd gen.) scalar leptoquarks simulated at four masses of 300 GeV, 400 GeV, 600 GeV, and 800 GeV with the MC generator Pythia [7] at 14 TeV pp center-of-mass energy. The next to leading order (NLO) cross sections for the above simulated signal decrease with leptoquark mass from a few pb to a few fb with mass point of 400 GeV at $(2.24 \pm 0.38)$ pb.

$^1 = 1$ would mean that leptoquarks do not decay into quarks and neutrinos.
Table I Partial cross-section (pb) that remains after each selection criterion for the 1st generation leptoquark channel. Baseline selection corresponds to the 1st selection of two electron candidates and two jets from the reconstructed objects. The Z/DY cross section is for the region M (ee) 60 GeV. VB pairs corresponds to the diboson processes of W W , W Z, and ZZ.

| Physics | Before Baseline | S_T | M ee | M_{1j} - M_{1j} mass window |
|---------|----------------|-----|------|------------------------------|
| LQ (400 GeV) | 2.24 | 1.12 | 1.07 | 1.00 | 0.534 |
| Z/DY 60 GeV | 1808 | 49.77 | 0.722 | 0.0664 | 0.0036 |
| tt | 450 | 3.23 | 0.298 | 0.025 | 0.0144 |
| VB pairs | 60.94 | 0.583 | 0.0154 | 0.0036 | 0.00048 |
| Multi-jet | 10^3 | 20.51 | 0.229 | 0.184 | 0.0 |

Table II Partial cross-section (pb) that remains after each selection criterion for the second generation leptoquark channel. Baseline selection corresponds to the 1st selection of two muon candidates and two jets from the reconstructed objects. The Z/DY cross section is for the region M ( ) 60 GeV. VB pairs corresponds to the diboson processes of W W , W Z, and ZZ.

| Physics | Before Baseline | p_T^{jet} | S_T | M ( ) | M_{1j} mass window |
|---------|----------------|----------|-----|------|-------------------|
| LQ (400 GeV) | 2.24 | 1.70 | 1.53 | 1.27 | 1.23 | 0.974 |
| Z/DY 60 GeV | 1808 | 79.99 | 2.975 | 0.338 | 0.0611 | 0.021 |
| tt | 450 | 4.17 | 0.698 | 0.0791 | 0.0758 | 0.0271 |
| VB pairs | 60.94 | 0.824 | 0.0628 | 0.00846 | 0.00308 | 0.00205 |
| Multi-jet | 10^3 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 |

2.1. Reconstruction and objects selection

Signal reconstruction requires selection of two high quality leptons and at least two jets. Each signal jet and lepton candidate is required to have transverse mass m_{T} > 20 GeV. This helps to suppress low p_T background predicted by the SM. Leptons are required to have pseudorapidity j j below 2.5, which is the inner detector’s acceptance, whereas jets are restricted to j j < 45 to suppress backgrounds from underlying event and minimum bias events that dominate in the forward region of the detector. In addition, leptons are required to pass lepton identification criteria, which, in case of electrons, are based on electron magnetic-shower shape variables in the calorimeter and, in the case of muons, are based on finding a muon track in the muon spectrometer and the inner detector together with a muon isolation requirement in the calorimeter. Electron candidates are also required to have a matching track in the inner detector. Furthermore, it is required that signal jet candidates are spatially separated from energy clusters in the electromagnetic calorimeter that satisfy electron identification criteria. Finally, a pair of leptoquark candidates are reconstructed from lepton-jet combinations. Given the fact that these four objects can be combined to give two pairs, the pair that has minimum mass difference between the two leptoquark candidates is assumed to be the signal.

2.2. Background Studies

The main backgrounds to the signal come from tt and Z+ jets production processes. Multi-jet production where two jets are misidentified as electrons, represents another background to the dielectron (1st gen.) channel. In addition, minor contributions arise from diboson production. Other potential background sources, such as single-top production, were also studied and found to be insignificant.

The backgrounds are suppressed and the signal significance is improved by taking advantage of the fact that the real state particles in signal-like events have relatively large p_T. A scalar sum of transverse momenta of signal jets and lepton candidates, denoted by S_T, helps in reducing the backgrounds while retaining most of the signal. The other variable used to increase the signal significance is the invariant mass of...
the two leptons, $M_{ll}$. The distributions of these two variables for the first generation channel are shown in Fig. 2.

![Figure 2: $S_T$ (top) and $M_{ll}$ of the selected electron pair after $S_T$ requirement (bottom) in 1st gen. leptoquark MC events ($m_{LQ} = 400$ GeV). Both distributions are normalized to 100 pb$^{-1}$ of integrated pp luminosity.](image)

After applying optimized selection on these two variables, $S_T$ and $M_{ll}$, relative contributions from the background processes from $t\bar{t}$, $Z+DY$, diboson and multijet are 22%, 7%, 0.4% and 18%, respectively. Partial cross-section for the signal and the background processes passing the selection criteria are shown in tables I and II for the 1st and 2nd generation channels, respectively. Figure 3 shows the invariant mass$^3$ of the reconstructed leptoquark candidates before and after background suppression criteria are applied to the MC data.

![Figure 3: Reconstructed electron-jet invariant mass for 1st gen. leptoquark ($m_{LQ} = 400$ GeV) in signal and background MC events after baseline selection (top) and after additional selection criteria based on $S_T$ and $M_{ll}$ (bottom) have been applied. Both distributions are normalized to 100 pb$^{-1}$ of integrated pp luminosity.](image)

2.3. Sensitivity and Discovery Potential

ATLAS’s sensitivity to leptoquark signal for a 400 GeV mass hypothesis and with an integrated pp luminosity of 100 pb$^{-1}$ is summarized in Fig. 4. The cross-sections include systematic uncertainties of 50%. Leptoquark-like events in the ATLAS detector are triggered by single leptons with an efficiency of 97%. ATLAS is sensitive to leptoquark masses of about 565 GeV and 575 GeV for 1st and 2nd generations, respectively, at the given luminosity of 100 pb$^{-1}$ provided the predicted cross-sections for the pair production of leptoquarks are correct.

3. Search for $W_R$ bosons and heavy Majorana neutrinos

$W_R$ bosons are the right-handed counterpart of the SM $W$ bosons. These right-handed intermediate vector bosons are predicted in LRSMs and can be produced at the LHC in the same processes as the SM’s $W$ and $Z$. They decay into heavy Majorana neutrinos. The Feynman diagram for $W_R$ production and subsequent decay to Majorana neutrino is shown in Fig. 4.

This section describes the analysis of $W_R$ production and its decays $W_R \rightarrow eN_e$ and $W_R \rightarrow N \rightarrow \ell\nu\ell\nu$. The...
Table III LRSM dilepton analysis. Partial cross-section (pb) that remains after each selection criterion for the dilepton channel.

| Physics sample | Before selection | Baseline selection M (jj) | M (eejj) | M (ee) | S_t |
|----------------|------------------|---------------------------|----------|--------|-----|
| LRSM 18.3     | 0.248            | 0.0882                    | 0.0861   | 0.0828 | 0.0786 |
| LRSM 15.5     | 0.470            | 0.220                     | 0.215    | 0.196  | 0.184 |
| Z / D Y 60 GeV| 1808.            | 49.77                     | 43.36    | 0.801  | 0.0132 |
| tt             | 450.             | 3.23                      | 3.13     | 0.215  | 0.0422 |
| V B pairs      | 60.94            | 0.583                     | 0.522    | 0.0160 | 0.0016 |
| Mu l i j t     | 10^3             | 20.51                     | 19.67    | 0.0490 | 0.0444 |

Table IV LRSM dimuon analysis. Partial cross-section (pb) that remains after each selection criterion for the dimuon channel.

| Physics sample | Before selection | Baseline selection M (jj) | M (eejj) | M (ee) | S_t |
|----------------|------------------|---------------------------|----------|--------|-----|
| LRSM 18.3     | 0.248            | 0.145                     | 0.141    | 0.136  | 0.128 |
| LRSM 15.5     | 0.470            | 0.328                     | 0.319    | 0.295  | 0.274 |
| Z / D Y 60 GeV| 1808.            | 79.99                     | 69.13    | 1.46   | 0.0231 |
| tt             | 450.             | 4.17                      | 4.11     | 0.275  | 0.0527 |
| V B pairs      | 60.94            | 0.824                     | 0.775    | 0.0242 | 0.0044 |
| Mu l i j t     | 10^3             | 0.0                       | 0.0      | 0.0    | 0.0 |

lowed by the decays N_e! e!q B N! q!q, which are detected in final states with (at least) two leptons and two jets. The two leptons can be of either same-sign or opposite-sign charge due to the M a j o r a n a nature of neutrinos. This analysis in both the dilepton and the dimuon channels has been performed without separating dileptons into same-sign and opposite-sign sam ples.

Studies of the discovery potential for W_R and M a j o r a n a neutrinos N_e and N! have been performed using MC samples where M (N_i) = 300 GeV; M (W_R) = 1800 GeV (referred to as LRSM 18.3) and M (N_i) = 500 GeV; M (W_R) = 1500 GeV (referred to as LRSM 15.5), simulated with PYTHIA according to a particular implementation of LRSM. The production cross-sections (pp(14 TeV) ! W_R X) times the branching fractions (W_R ! N_e! e!q B N! q!q) are 24.8 pb and 47 pb for LRSM 18.3 and LRSM 15.5, respectively.

3.1. Reconstruction and objects selection

Signal event candidates are reconstructed using two electron or muon candidates and two jets that pass the standard selection criteria as discussed in section 3.2. The two signal jet candidates are combined with each of the signal leptons and the combination that gives the smallest invariant mass is assumed to be the new heavy neutrino candidate. The other remaining lepton is assumed to come directly from the decay of the W_R boson. If signal electrons and signal jets overlap in R within 0.4 then, to avoid double counting, only the two signal jets are used to reconstruct the invariant masses of the heavy neutrino candidate and W_R.

3.2. Background Studies

The main backgrounds to the LRSM analyses studied here are the same as mentioned in section 3.2. The same background suppression criteria as in the lepton-quark analyses are also effective here, namely S_T and m_T. The distributions of these two variables for the dilepton channel are shown in Fig. 8a. Partial cross-section for the signal and the background processes passing the selection criteria are shown in tables 11 and 12 for the dilepton and dimuon channels, respectively. Figure 8b shows the invariant mass of the reconstructed W_R candidates before and after background suppression criteria are applied to the MC data.

3.3. Sensitivity and Discovery Potential

Signal significance for W_R analyses in the dilepton and dimuon channels as a function of integrated pp luminosity at 14 TeV is summarized in Fig. 8c. The results include systematic uncertainty of 45% and 40% for di-lepton and dimuon channel, respectively.
events in this analysis are also triggered by single leptons with an efficiency of 97%.

4. Conclusions

Dilepton–jet based final states have been discussed in both electron and muon channels. Discovery potential for leptoquarks and LRSM with early LHC data have been investigated with the predicted cross-sections for these models. Assuming $\alpha = 1$, both 1st and 2nd generations leptoquarks could be discovered with masses up to 550 GeV with 100 pb$^{-1}$ of data. Two LRSM mass points LRSM$_{18.3}$ and LRSM$_{15.5}$ for the $W_R$ bosons and heavy Majorana neutrinos have been studied. The discovery of these new particles with such masses would require integrated luminosities of 150 pb$^{-1}$ and 40 pb$^{-1}$, respectively.
Figure 7: LRSM analysis. The distributions of the reconstructed invariant masses for $W_R \rightarrow N$ candidates in background and signal (LRSM $18.3$ and LRSM $15.5$) events before (top) and after (bottom) background suppression is performed in dimuon channel analysis. Both distributions are normalized to $100 \text{ pb}^{-1}$ of integrated pp luminosity.

Figure 8: LRSM analysis. Expected signal significances versus integrated pp luminosity for $N$, $N$ and $W_R$ mass hypotheses, according to signal MC samples LRSM $18.3$ and LRSM $15.5$. Open symbols show sensitivities without systematic uncertainties. Sensitivities shown with closed symbols include an overall relative uncertainty of $45\%$ ($40\%$), estimated for background contributions in the dielectron (dimuon) analysis.

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