Multi-frequency GMRT Observations of the H II regions S 201, S 206, and S 209

Galactic Temperature Gradient

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Abstract. We present radio continuum images of three Galactic H II regions, S 201, S 206, and S 209 near 232, 327, and 610 MHz using the Giant Meterwave Radio Telescope (GMRT). The GMRT has a mix of short and long baselines, therefore, even though the data have high spatial resolution, the maps are still sensitive to diffuse extended emission. We find that all three H II regions have bright cores surrounded by diffuse envelopes. We use the high resolution afforded by the data to estimate the electron temperatures and emission measures of the compact cores of these H II regions. Our estimates of electron temperatures are consistent with a linear increase of electron temperature with Galacto-centric distance for distances up to ∼ 18 kpc (the distance to the most distant H II region in our sample).

Key words. H II regions – ISM: individual – S 201, S 206, S 209: radio continuum – low frequency

1. Introduction

A number of studies have indicated that the electron temperature $T_e$ of H II regions increases with increasing Galacto-centric distance (e.g. Deharveng et al. 2000 and references therein). This effect is attributed to a decrease in heavy elements abundances with Galacto-centric distance. A low metal abundance leads to less effective cooling and consequently higher electron temperature. These studies are based either on estimates of $T_e$ from radio recombination lines (RRLs) (which in turn depend on corrections for departures from local thermodynamic equilibrium (LTE) and for collisional broadening effects), or estimates based on line strengths of the forbidden line transitions of Oxygen [O III] $\lambda\lambda$4363, 5007 (which are strongly dependent on temperature variations, if any, over the observed volume). Further, most of these studies are based on observations of H II regions with Galacto-centric distances $R_G \leq 15$ kpc with very few measurements of $T_e$ beyond 15 kpc. Consequently most determinations of metallicities of the outer galaxy H II regions are based on values of $T_e$ taken from an extrapolation of the observed gradient in temperature up to about 15 kpc (e.g., Deharveng et al. 2000). Since the O/H ratio (a commonly used indicator of metal abundance) depends sensitively on $T_e$, metallicities of the outer galaxy H II regions are poorly constrained. In view of this, it is important to get independent estimates of the electron temperatures of H II regions in the outer galaxy.

An independent measurement of $T_e$ can be obtained from radio continuum observations. The ionized material in H II regions emits radio continuum through free-free emission. At sufficiently low radio frequencies where the nebula is optically thick ($\tau \gg 1$), the emergent radiation is a black body spectrum, and therefore, the observed brightness temperature is equal to the electron temperature $T_e$. On the other hand, at sufficiently high radio frequencies, where the optical depth $\tau$ of thermal electrons is low ($\tau \ll 1$), the observed brightness is proportional to the emission measure of the nebula. Most of the available radio maps for H II regions are at high radio frequencies (i.e. above 1.4 GHz, e.g., Fich 1993, Balser et al. 1995). These maps show that H II regions often have a bright core with several knots surrounded by an extended envelop of diffuse emission. These core–envelope structures of H II regions imply that accurate measurement of $T_e$ from low radio frequency observations requires high angular...
resonation, since, often only bright compact cores will be optically thick at frequencies of a few hundred MHz. This study presents an analysis of the low-frequency GMRT observations of three Galactic diffuse H II regions spanning Galacto-centric distances up to 18 kpc.

The GMRT is an ideal telescope for these observations since it operates at several low radio frequency bands, viz., 150, 232, 327, 610, and 1420 MHz and also it has a hybrid configuration which makes it sensitive to both diffuse emission (on scales up to ~ 45’ at 232, 30’ at 327, and 17’ at 610 MHz) while also having the resolution (~ 15” at 232, 10” at 327, and 6” at 610 MHz) to resolve the compact cores.

2. Observations

The observations were carried out during the period of August to December, 1999 at three frequency bands, viz., 232, 327, and 610 MHz. The GMRT has a ‘Y’ shaped hybrid configuration of antennas with six antennas along each of the three arms and twelve antennas randomly placed in a compact arrangement near the centre of ‘Y’ (for details, see Swarup et al. 1991). The compact array at the centre is about a kilometer across and is generally referred as the “central square”. Baselines in the central twelve antennas randomly placed in a compact arrangement near the centre of ‘Y’ (for details, see Swarup et al. 1991). The observations were carried out with typically 20 to 25 antennas in different observing sessions.

The data were recorded in the default correlator mode which produces visibilities in 128 channels over a user selectable bandwidth in multiples of 2 starting from 62.5 kHz and up to 16 MHz. The observational parameters are summarized in table 1.

The observations near 610 and 327 MHz were made using the full 16 MHz bandwidth while observations near 232 MHz were made with a bandwidth of 2 MHz centered at a frequency around which least local interference has been detected in the past observations. The images at all frequencies are however made using data from only one channel which corresponds to a bandwidth of 125 kHz at 327 and 610 MHz, and 15.6 kHz at 232 MHz. This restriction was partly because of a crunch in disk storage at the time when these data were taken, and partly because dynamic range limitations at the GMRT at the time we took the data meant that the increase in bandwidth did not result in a proportionate increase in sensitivity. At each frequency band, we observed the source for about 8–10 hours, primarily in order to have a good (u, v) coverage.

For all the observations, the source 3C 48 was used as the primary flux calibrator. The flux density of 3C 48 at each frequency was estimated using the Baars et al. (1977) flux densities of standard VLA calibrators. The phase and amplitude gains of antennas were derived from observations of a secondary calibrator at intervals of 45 minutes. For observations on S 206 and S 209, 3C 119 and 0107+562 were used as a secondary calibrator while 0107+562 was used as a secondary calibrator for observations on S 201. Both 3C 119 and 0107+562 are standard VLA calibrators. The fluxes of secondary calibrators were determined via boot-strapping the fluxes of the primary calibrator 3C 48.

The data were carefully checked for interference or other problems. At 232 and 327 MHz, a few short baselines were found to be corrupted, possibly by interference, and were removed. The data at 610 MHz were found to be free from any interference. Data reduction was done in classic AIPS. The calibrated data were Fourier transformed using appropriate (u, v) ranges, tapers and weights to make different images, some of which are sensitive to large scale structures, and others which have the maximum possible angular resolution. These images were deconvolved using the ‘CLEAN’ algorithm as implemented in AIPS task ‘IMAGR’. The final gains of the antennas were fixed using several iterations of self-calibration.

The variations in system temperatures of GMRT antennas are currently not routinely monitored during observations. The system temperature at 610 MHz was measured both toward the absolute flux calibrator 3C 48 and the target source by firing the noise calibration diodes. For 327 and 232 MHz images, the system temperature toward 3C 48 and target source were obtained using interpolated values of sky temperature from 408 MHz all–sky map of Haslam et al. (1982). A correction factor equal to the ratio of the system temperature toward the target source and 3C 48 was applied in the deconvolved image. The deconvolved images were finally corrected for the primary beam attenuation, assuming a Gaussian shape for the primary beam. The half power

### Table 1. Observational Details

| Field Centre (B1950) RA | Field Centre (B1950) Dec | Frequency (MHz) | Duration of observation (Hours) | Range of baselines (k.l) | rms noise in the image (mJy beam⁻¹) |
|------------------------|--------------------------|-----------------|---------------------------------|--------------------------|-----------------------------------|
| 02°59′12″                | 60°17′00″                | 231             | 8                               | 0.05–15                  | 2.5                               |
|                        |                          | 616             | 8                               | 0.09–25                  | 1.2                               |
| 03°59′24″                | 51°11′00″                | 236             | 9                               | 0.05–18                  | 7.4                               |
|                        |                          | 328             | 10                             | 0.06–27                  | 3.0                               |
| 04°07′18″                | 51°02′00″                | 328             | 9                               | 0.10–26                  | 2.0                               |
|                        |                          | 613             | 10                             | 0.10–50                  | 1.0                               |

The observations near 610 and 327 MHz were made using the full 16 MHz bandwidth while observations near 232 MHz were made with a bandwidth of 2 MHz centered at a frequency around which least local interference has been detected in the past observations. The images at all frequencies are however made using data from only one channel which corresponds to a bandwidth of 125 kHz at 327 and 610 MHz, and 15.6 kHz at 232 MHz. This restriction was partly because of a crunch in disk storage at the time when these data were taken, and partly because dynamic range limitations at the GMRT at the time we took the data meant that the increase in bandwidth did not result in a proportionate increase in sensitivity. At each frequency band, we observed the source for about 8–10 hours, primarily in order to have a good (u, v) coverage.

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can be modeled by solving equations (1) and (2) iteratively for \( d \) (using table 6 of Mezger & Henderson 1967) for the frequency range 200–600 MHz and frequency and optically thin at the other.

when observations at least two frequencies are available and the frequencies are such that the H\( \alpha \) is a correction factor which depends both upon the temperature and frequency. We have used an average value of this case, since the cores are unresolved, is taken to be the synthesized beam size), and

The emission measure

1.4 GHz VLA image of Fich (1993). The high resolution 231 MHz GMRT image (Fig. 1\[B\]) shows the core to be a complex structure consisting of several unresolved compact sources. The di

616 MHz GMRT low resolution image (Fig. 1\[C\]) traces di

15 GHz radio continuum image reveals a bright arc like core with multiple peaks of emission (Felli et al. 1987).

The 616 MHz GMRT low resolution image (Fig. 1\[C\]) traces diffuse emission extending up to ~5′ which is consistent with the 1.4 GHz VLA image of Fich (1993). The high resolution 231 MHz GMRT image (Fig. 1\[B\]) shows the core to be a complex structure consisting of several unresolved compact sources. The diffuse nebulosity extending toward the west of the core in the 231 MHz image (Fig. 1\[B\]) is consistent with the 15 GHz radio image of Felli et al. (1987).

Images of S 206 are shown in Fig. 1\[D\], [E], & [F]. S 206, \((l = 150.74, b = -0.75;\) also known as NGC 1491), is an evolved H\( \pi \) region at a Galacto-centric distance of 11.1 kpc (Deharveng et al. 2000). The excitation is believed to be provided by a cluster of OB stars (Chini & Wink 1984). Our high resolution 613 MHz image (Fig. 1\[G\]) shows the core region to consist of an asymmetric, incomplete ring like structure. The high resolution image at 328 MHz (not shown here) is morphologically very similar to the 613 MHz image. The 236 MHz image (Fig. 1\[H\]) while showing overall similarity to the 613 MHz and 328 MHz maps, does show some di

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4. Discussion

We use these low frequency images to estimate electron temperatures and emission measures of the compact cores of the H\( \pi \) regions. If we approximate these cores to be homogeneous and spherically symmetric, then the flux \( S \) is given by

\[
S = 3.07 \times 10^{-2} T_e \nu^2 \Omega (1 - e^{-\tau(\nu)})
\tag{1}
\]

\[
\tau(\nu) = 1.643 a \times 10^5 \nu^{-2.1} EM T_e^{-1.35}
\tag{2}
\]

(Mezger & Henderson 1967) where \( S \) is the integrated flux density in Jy, \( T_e \) is the electron temperature in Kelvin, \( \nu \) is the frequency of observation in MHz, \( \tau \) is the optical depth, \( \Omega \) is the solid angle subtended by the source in steradian, (which in this case, since the cores are unresolved, is taken to be the synthesized beam size), and \( EM \) is the emission measure in cm\(^{-6}\) pc. The emission measure \( EM \) is defined as \( \int n_e^2 dl \); the integral being taken along the line of sight and averaged over the beam. \( a \) is a correction factor which depends both upon the temperature and frequency. We have used an average value of \( a \) as 0.98 (using table 6 of Mezger & Henderson 1967) for the frequency range 200–600 MHz and \( T_e \sim 10000 \) K. The H\( \pi \) region cores can be modeled by solving equations (1) and (2) iteratively for different \( EM \) and \( T_e \). The fitting procedure converges rapidly when observations at least two frequencies are available and the frequencies are such that the H\( \pi \) region is optically thick at one frequency and optically thin at the other.

| Name | \( \alpha, \delta \) (1950) | Frequency | Flux | Area | \( T_e \) | \( EM \) |
|------|----------------|-----------|-----|------|--------|-------|
|      | \( h m s \)      | (MHz)     | (Jy)| (arcmin\(^2\)) | (K)    | (cm\(^{-6}\) pc) |
| S 201 | 02 59 20.1 60 16 10 | 231       | 0.78| 16   | 7070 ± 1100 | 1.02(±0.05) \times 10^5 |
|      | 616              |           | 1.15| 38   |         |       |
| S 206 | 03 59 24.0 51 11 00 | 236       | 16.3| 329  | 8350 ± 1600 | 3.93(±0.40) \times 10^5 |
|      | 328              |           | 18.2| 347  |         |       |
|      | 613              |           | 20.0| 350  |         |       |
| S 209 | 04 07 20.1 51 02 30 | 236       | 13.4| 267  | 10855 ± 3670 | 2.58(±0.29) \times 10^5 |
|      | 328              |           | 16.6| 372  |         |       |
|      | 613              |           | 17.0| 386  |         |       |

points (HPBW) of the primary beam of GMRT antenna are estimated as 1.85, 1.35, and 0.72 degree for 232, 327 and 610 MHz respectively.
We measured the peak flux densities of cores after convolving the images of a H\(\text{\textsc{n}}\) region at different frequencies to a common angular resolution (i.e. the source size \(\Omega\) was taken to be 1.133 \(\times \theta_a \times \theta_b\), where \(\theta_a\) and \(\theta_b\) are the half power points of the common convolved beam). The best fit values for \(T_e\) and \(EM\) as obtained from the fitting procedure described above are listed in table 3 and the observed and model fluxes are plotted in Fig. 2. The columns in table 2 are as follows. col. (1): Name of the H\(\text{\textsc{n}}\) region, col. (2): Coordinates (right ascension, declination) of the core for which the electron temperature has been measured, col. (3): The frequency of observation, col. (4): Integrated flux of the entire H\(\text{\textsc{n}}\) region, col. (5): Area over which radio emission is detected.

**Fig. 1.** [A] S 201 at 616 MHz. The angular resolution is 26" \(\times\) 17". [B] S 201 at 231 MHz. The angular resolution is 15" \(\times\) 13". [C] S 201 at 616 MHz made using u-v range 0–1 k\(\lambda\) only. The angular resolution is 133" \(\times\) 129". [D] S 206 at 613 MHz. The angular resolution is 13" \(\times\) 11". The regions marked as A, B, C, & D are from the 5 GHz image of Deharveng et al. (1976). [E] S 206 at 236 MHz. The angular resolution is 20" \(\times\) 20". [F] 328 MHz image of S 206 made using u-v range 0–1 k\(\lambda\) only. The synthesized beam is 180" \(\times\) 149". [G] S 209 at 613 MHz. The angular resolution is 10" \(\times\) 10". [H] S 209 at 236 MHz. The angular resolution is 25" \(\times\) 25". [I] 613 MHz image of S 209 made using u-v range only up to 1 k\(\lambda\). The angular resolution is 160" \(\times\) 136".

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and over which the flux has been integrated to get the value listed in col. (4), col. (6): Estimated electron temperature of the core, col. (7): Estimated emission measure of the core.

The electron temperature of S 201 is estimated to be $7070 \pm 1100$ K toward the peak radio emission. The earlier estimate for $T_e$ toward S 201 was $\sim 5000$ K based on non-detection of [O iii] lines, 4959, 5007 (Mampaso et al. 1989). The electron temperature of 8350±1600 K, derived for the core of S 206 (knot–A in Fig. 1[C]) is in reasonable agreement with previous measurements, viz. 8400 ± 800 K obtained using the H94 α recombination line by Carral et al. (1981) and 9118 K obtained from the [O iii]4363, 5007 lines ratio (Deharveng et al. 2000). The emission measure is $3.93(\pm 0.40) \times 10^5$ cm$^{-6}$ pc, consistent with the value obtained by Deharveng et al. (1976). For S 209, the electron temperature corresponding to the peak radio emission at 613 MHz is estimated to be $10855 \pm 3670$ K, somewhat higher than the value of 8280 K obtained using the H137 β recombination line by Churchwell et al. (1978) but in reasonable agreement with the estimate of 11000 K which was derived from H91α & H114β recombination lines (Balser et al. 1994).

Figure 3 is a plot of the electron temperature vs. Galacto-centric distance for the three H II regions observed at the GMRT. The estimated emission measures and electron temperatures are listed in Table 2.

5. Conclusions

Three outer galaxy H II regions, S 201, S 206 and S 209 have been imaged at meter wavelengths using the GMRT. The images of these H II regions have been obtained at a resolution of less than a pc. This is the highest resolution achieved for any H II region at such low radio frequencies. All three H II regions show structures down to our resolution limit. The high resolution images near 610 MHz of these H II regions show a good correspondence with the radio continuum images at cm wavelengths. The low resolution radio images show that these H II regions are surrounded by large diffuse envelopes. The high resolution radio images have allowed us to get estimates of $T_e$ of these H II regions. From these measurements we find that:

1. the estimates of $T_e$ are in general consistent with that obtained from RRLs and [O iii] lines, 4363, 5007 line measurements, and
2. the measured temperatures are consistent with a linear increase of $T_e$ with Galacto-centric distance until $R_G \sim 18$ kpc.

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Fig. 3. The electron temperature vs. Galacto-centric distance for three H\textalpha{} regions studied in this paper. The solid line is the relationship derived by Deharveng et al. (2000) based on a sample of six H\textalpha{} regions spanning a Galacto-centric distance from 6.6 to 14.8 kpc.

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