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A reversible single-molecule switch based on activated antiaromaticity

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Single-molecule electronic devices provide researchers with an unprecedented ability to relate novel physical phenomena to molecular chemical structures. Typically, conjugated aromatic molecular backbones are relied upon to create electronic devices, where the aromaticity of the building blocks is used to enhance conductivity. We capitalize on the classical physical organic chemistry concept of Hückel antiaromaticity by demonstrating a single-molecule switch that exhibits low conductance in the neutral state and, upon electrochemical oxidation, reversibly switches to an antiaromatic high-conducting structure. We form single-molecule devices using the scanning tunneling microscope–based break-junction technique and observe an on/off ratio of ~70 for a thiophenylene derivative that switches to an antiaromatic state with 6-4-6-π electrons. Through supporting nuclear magnetic resonance measurements, we show that the doubly oxidized core has antiaromatic character and we use density functional theory calculations to rationalize the origin of the high-conductance state for the oxidized single-molecule junction. Together, our work demonstrates how the concept of antiaromaticity can be exploited to create single-molecule devices that are highly conducting.

RESULTS AND DISCUSSION

We first characterize the solution-based optical and electronic properties of the TBTP by cyclic voltammetry (CV) and spectroelectrochemistry. Figure 2B shows a CV plot where we see four reversible redox peaks, two corresponding to oxidation and two corresponding to reduction. We note that the two oxidation peaks at \( E_{\text{ox1}} = 0.41 \) V and \( E_{\text{ox2}} = 0.56 \) V (versus ferrocene [Fc\(^{+}\)]) are close in energy, whereas the reduction peaks are well separated. To characterize the changes in optical properties of TBTP under various applied electrochemical potentials, we measure ultraviolet–visible (UV-vis) absorption spectra of TBTP at different potentials. We start with the neutral molecule that exhibits an optical energy gap of 1.77 eV (fig. S1). Upon applying a potential of 0.9 V (versus ferrocene [Fc\(^{+}\)]) the molecule is doubly oxidized (+2 state), and the optical gap decreases to about 0.85 eV (Fig. 2C). The spectroelectrochemical changes are reversible, as is visible from the CV and additional UV-vis spectra (fig. S2), which shows the spectral characteristics of the transient radical cation. To probe changes to the molecular structure after oxidation, we chemically oxidize TBTP with AgSbF\(_5\)/I\(_2\) and this shows further evidence of the +2 state by matching the UV-vis spectra to the spectroelectrochemistry data (the radical cation is not observed by chemical oxidation). TBTP can be chemically oxidized reversibly (fig. S3), and the dication exhibits remarkable stability at ambient conditions for several hours. Considering that the antiaromaticity of this stable system can be probed by proton nuclear magnetic resonance (\(^{1}H\)-NMR) (23), we collected the \(^{1}H\)-NMR spectra for the neutral and chemically oxidized TBTP at room temperature. The NMR spectra show that the proton signal from the benzene ring of the BTP core exhibits a marked upshielding shift, indicating antiaromatic character (fig. S4). Although this upshielding shift is not seen for all protons on the core due to charge delocalization, this key result is
consistent with the NMR spectra predicted by density functional theory (DFT) calculations, as detailed in Materials and Methods. Thus, it can be concluded that the BTP core takes on antiaromatic character in the +2 state, consistent with the simple electron-counting picture presented above.

To test the hypothesis that antiaromaticity can lead to highly conductive molecular devices, we performed single-molecule measurements under an electrochemical potential using the scanning tunneling microscope break-junction (STM-BJ) technique (24, 25). The details of the method have been described before by Capozzi et al. (24). Briefly, we use a gold STM tip and substrate, which forms the working electrodes, whereas an additional Pt wire is introduced into the fluid cell serves as the counter (or gate) electrode. Measurements are carried out in 0.1 mM TBTP solution in propylene carbonate (PC), a polar solvent, with 0.1 M tetrabutylammonium perchlorate as the supporting electrolyte, and potential calibrations are carried out by adding ferrocene to the solution and measuring its reversible oxidation potential. The tip is insulated with Apiezon wax to reduce the background ionic current (26). The conductance of TBTP is first measured without applying a gate voltage. Figure 2D shows the one-dimensional (1D) conductance histograms of TBTP at a bias of 0.09, 0.45, and −0.45 V relative to the substrate. Each conductance histogram reveals a clear narrow peak at a bias-dependent conductance value. The inset of Fig. 2D shows the peak conductance plotted against the tip bias (see fig. S5 for full data sets), where a modest conductance increase of a factor of 5 is seen within this bias range. Such a bias polarity-dependent conductance indicates that the highest occupied molecular orbital (HOMO) dominates charge transport (27).

The electrochemical switching of the single-molecule junction was accomplished by changing the voltage applied to the Pt gate electrode (relative to the substrate) and repeating the measurements while maintaining a low bias voltage of 25 mV on the tip (also relative to the substrate). Note that in this STM-BJ measurement, the voltage is applied to the gate electrode, relative to the substrate. Therefore, oxidation events are observed at negative gate voltages. This is corroborated by the in situ electrochemistry shown in fig. S6A. By contrast, in traditional CV measurements, the potential is applied to the working electrode, relative to the reference. Conductance histograms measured at different voltages applied to the gate are shown in Fig. 3A. At the positive gate voltage, the conductance peak shifts to lower values, consistent with transport being dominated by the HOMO. At negative gate voltages, a modest increase of the conductance is first observed until an applied gate voltage of about −1.4 V (+1 V measured potential relative to Fc+/0). Beyond −1.4 V, the conductance peak shifts markedly by a factor of ~70 (Fig. 3A), and the histograms broaden—an aspect we will discuss further below. This sharp increase in conductance is visible in the inset of Fig. 3A, where we plot the measured conductance against the gate potential (see fig. S5B for full data sets). We show further that the switching is reversible by changing the gate potential between negative (<−1.5 V) and positive (>1 V) values, in an alternating fashion, after every 100 conductance traces. The conductance peak values determined from the histogram of these traces (fig. S6B) are plotted in Fig. 3B, yielding an average on/off ratio value of ~70. Note that the redox-switching experiments were conducted at ambient conditions, without the rigorous exclusion of air and moisture. In addition, to show that the gate potential does change the oxidation state of the molecules in solution in our STM setup, we record the in situ linear sweep voltammetry (LSV), measuring the current through the STM tip as a function of the applied gate voltage. We do this measurement after adding ferrocene to the solution to calibrate the applied gate voltage and confirm oxidation states (fig. S6A). We see a broad oxidation peak at a gate voltage of −1.4 V, consistent with our finding that a negative gate voltage needs to be applied to measure the conductance of the oxidized TBTP.

To confirm that we are measuring the conductance of a TBTP molecule bridging the gap between two Au electrodes, we compare 2D conductance-displacement histograms compiled from these data in Fig. 3 (C and D) (28). Both figures show a molecular feature that extends by about 1 nm relative to the point where the gold point contact is ruptured. We have previously shown that the extent of the 2D histogram feature relates to the molecular backbone length (28), and thus, these data reveal that a stable molecular junction is formed at the high negative gate bias. The key result of this study is that changing the oxidation state of TBTP without making/breaking covalent bonds results
in a large change in conductance, which can be attributed to the increased antiaromatic character of the oxidized species.

To provide further support for the observed increased conductance upon oxidation, we turn to DFT-based calculations of the single-molecule junction. We first construct the molecular junction containing the neutral TBTP (see Fig. 4A for the geometry) and perform conductance calculations with a DFT-based nonequilibrium Green’s function (NEGF) approach, with the TranSIESTA package (see Materials and Methods for details) (29). Figure 4B shows the resulting transmission as a function of energy (orange trace). The conductance at zero bias (transmission at Fermi energy) is small and corresponds to a junction in its “off” state. We note that for a quantitative comparison between theory and experiment, self-energy corrections to the DFT mean-field eigensystem are required, as in, for example, the approximate DFT + Σ method (see Materials and Methods and the Supplementary Materials) (30). Here, we focus on the comparison between the neutral and the oxidized TBTP molecule. To model the molecule in a +2 oxidation state, eight fluorine atoms are added around the BTP core at approximately 1 Å above and below the central rings, such that the total Mulliken charge on the molecule reaches +2 in the gas phase (see Fig. 4A for the geometry). We then perform a DFT-NEGF calculation of the molecular junction with the fluorine atoms present. The resulting transmission is shown in Fig. 4B.

The transmission functions of both the neutral and oxidized species show stark differences. The oxidized TBTP leads to a partially occupied resonance and much higher transmission close to the Fermi energy. This increase corresponds to a conductance enhancement of a factor of ~60, consistent with the experimental results, where the average on/off ratio is ~70. Although the actual conductance of the oxidized species may be sensitive to the parameters of the calculation, the qualitative trend of increasing conductance after oxidation is correctly captured here. Because the resonance close to the Fermi level is quite narrow, small changes in the metal work function as a result of changes in junction geometry or ionic environment could result in large changes in conductance. This is consistent with the experimental data that shows a broader conductance peak for the +2 state of TBTP.

**Fig. 4. DFT-based transmission calculations.** (A) The molecular junction containing the neutral TBTP molecule (H, light pink; C, brown; and S, yellow) and eight fluorine atoms at about 1 Å above (dark blue) and below (light blue) the central rings to model the TBTP molecule in a +2 oxidized state. (B) Transmission as a function of energy for the neutral (orange) and the oxidized molecule (green). $E_F$ is the Fermi level of the junction.

**MATERIALS AND METHODS**

**Molecular characterizations**

$^1$H-NMR and $^{13}$C-NMR were recorded on a Bruker Avance III 500 (500 MHz) spectrometer in dichloromethane-$d_2$ [residual solvent peak at $\delta = 5.32$ parts per million (ppm) for $^1$H NMR] and chloroform-$d$ solution (residual solvent peak at $\delta = 77.16$ ppm for $^{13}$C NMR). Mass spectra were obtained at the Columbia University Mass Spectrometry Facility using a Xevo G2-XS (Waters) equipped with a quadrupole time-of-flight detector with multiple inlet and ionization capabilities, including electrospray ionization, atmospheric pressure chemical ionization, and atmospheric solids analysis probe. The base peaks were usually obtained as [M]$^+$ or [M + H]$^+$ ions. UV-vis absorption data were acquired on a Varian Cary 5000 UV-Vis-NIR spectrophotometer. Cyclic voltammograms were recorded on a CH Instruments 660 electrochemical workstation using Pt plate electrode as the working electrode, Pt wire as the counter electrode, and Ag/AgCl electrode as the reference electrode at room temperature.

**STM-BJ measurements**

Conductance measurements were performed using a custom STM that has been described in detail before (31). Conductance measurements were performed in dilute solutions (10 to 100 μM) in PC with tetrabutylammonium perchlorate as the supporting electrolyte. The insulated tips were created by driving a mechanically cut gold tip through Apiezon wax (26). 1D conductance histograms were constructed using...
logarithmic bins (100 bins/decade), and 2D histograms used logarithmic bins along the conductance axis (100 bins/decade) and linear bins along the displacement axis. All histograms were constructed without any data selection.

**Antiaromaticity characterizations**

To characterize the antiaromatic character of the oxidized TBTP\textsuperscript{3+}, the NMR spectra of the neutral and oxidized species were collected at room temperature (23). The oxidant, AgSbF\textsubscript{6}/I\textsubscript{2}, was added to the solution of TBTP in deuterated dichloromethane (CD\textsubscript{2}Cl\textsubscript{2}) in a 2:1 ratio (oxidant/TBTP) (32).

**Calculation of the NMR spectra and nuclear-independent chemical shifts**

We calculated the NMR spectra and nuclear-independent chemical shift (NICS) using DFT within the Gaussian 09 suite (33). All geometries were optimized with B3LYP functional and the 6-31G** basis set. NMR and NICS calculations were conducted with the GIAO-B3LYP/6-31G** method. The calculated NMR chemical shifts are generally consistent with the experimental ones as indicated in table S1. The change of aromaticity for all the rings in the BTP core can be determined from NICS calculations. In general, a positive NICS value indicates an antiaromatic character, whereas a negative NICS value indicates aromatic character (34, 35). We calculated two types of NICS values: NICS(1)zz, which represents the NICS 1 Å above and below center mass of the ring, and NICS(0), which gives NICS value at the center of mass of the ring. Results are reported in table S2.

**Transport calculations**

We relaxed the junction geometry using the Perdew-Burke-Ernzerhof (PBE) functional (36), as implemented in SIESTA package (37). Pseudo-potentials and basis sets were adapted from a previous work (30). Each electrode consisted of seven layers of 36 gold atoms forming a 6 × 6 unit cell of the Au(111) surface. In addition, the S atoms in the molecule were attached to the gold substrate using trimer binding motifs. During geometry relaxation, the outer four layers of gold on each side were treated as a "rigid body" and were only allowed to relax along the transport (z) direction, as x and y coordinates were held fixed and the internal atomic distances within this set of gold atoms were kept at bulk values. All x, y, and z coordinates of the molecule, the trimer tip motif, and the inner three layers of gold atoms on each side were allowed to fully relax. In this way, we achieved the optimal geometry around the molecule-gold contact while maintaining the bulk gold geometry far away from the molecule. Periodic boundary conditions with large vacuum along the z direction and a 3 × 3 × 1 \( \mathbf{k} \)-mesh were used in the geometry relaxation, which was performed until all forces were <0.04 eV/Å. After the junction was relaxed, transport properties were calculated using the TranSIESTA package (29) with the same functional, pseudopotentials, basis sets, and \( \mathbf{k} \)-mesh as above.

It is known that the standard DFT-NEGF formalism overestimates the zero-bias conductance due to underestimation of the level alignment between frontier orbitals of the molecule and the Fermi level of the junction (30), especially when local or semilocal functionals are used. We used the DFT + \( \Sigma \) method (30, 38) to correct the level alignment and transport properties. The correction to level alignment consists of two terms: a self-energy correction to the gas-phase HOMO-LUMO (lowest unoccupied molecular orbital) gap and a correction to the electron addition/removal energies due to the electrostatic polarization (image charge) of the electrodes and environment. For the gas-phase correction, we used the optimally tuned range-separated hybrid functional (39) to accurately calculate the HOMO energy of the gas-phase molecule. For the TBTP molecule, it is \(-6.2\) eV (optimal Coulomb separation parameter is 0.11 bohr\textsuperscript{-1}), yielding a gas-phase correction to the PBE HOMO of 1.8 eV downward. For the "image-charge" contribution, we used the image plane of 0.9 Å, as determined previously (40), which results in an image-charge correction of 0.5 eV. The total self-energy correction for HOMO resonance is then 1.3 eV downward. To model the oxidized molecule, we decorated the region around the BTP core with eight fluorine atoms and repeated the transmission calculations. However, we cannot use the DFT + \( \Sigma \) method to calculate the transmission for the molecule-fluorine system due to a strong charge transfer between the molecule and the fluorine atoms.

**SUPPLEMENTARY MATERIALS**

Supplementary material for this article is available at http://advances.sciencemag.org/cgi/content/full/3/10/eaao2615/DC1

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