Millimeter Imaging of MWC 758: Probing the Disk Structure and Kinematics

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ABSTRACT

We investigate the structure and kinematics of the circumstellar disk around the Herbig Ae star MWC 758 using high-resolution observations of the 12CO (3–2) and dust continuum emission at the wavelengths of 0.87 and 3.3 mm.

We find that the dust emission peaks at an orbital radius of about 100 AU, while the CO intensity has a central peak coincident with the position of the star. The CO emission is in agreement with a disk in Keplerian rotation around a 2.0 $M_\odot$ star, confirming that MWC 758 is indeed an intermediate-mass star. By comparing the observation with theoretical disk models, we derive that the disk surface density $\Sigma(r)$ steeply increases from 40 to 100 AU and decreases exponentially outward. Within 40 AU, the disk has to be optically thin in the continuum emission at millimeter wavelengths to explain the observed dust morphology, though our observations lack the angular resolution and sensitivity required to constrain the surface density on these spatial scales. The surface density distribution in MWC 758 disk is similar to that of “transition” disks, though no disk clearing has been previously inferred from the analysis of the spectral energy distribution (SED). Moreover, the asymmetries observed in the dust and CO emission suggest that the disk may be gravitationally perturbed by a low-mass companion orbiting within a radius of 30 AU. Our results emphasize that SEDs alone do not provide a complete picture of disk structure and that high-resolution millimeter-wave images are essential to reveal the structure of the cool disk mid-plane.

Key words: planet–disk interactions – protoplanetary disks – submillimeter: planetary systems – techniques: interferometric

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1. INTRODUCTION

Millimeter-wave observations of circumstellar disks provide the best tool to measure the distribution and kinematics of the circumstellar material. These observations have for years achieved 1″–3″ angular resolution to resolve the disk emission on scales down to an orbital radius of about ~70 AU at the distance of the nearby star-forming regions (e.g., Koerner & Sargent 1995; Guilloteau & Dutrey 1998; Kitamura et al. 2002). More recent observations are reaching subarcsecond resolution that is resolving disks on scales of the Kuiper Belt and yielding unprecedented information on the innermost disk regions (see, e.g., Andrews et al. 2009; Brown et al. 2009; Hughes et al. 2009; Isella et al. 2009, 2010; Piétu et al. 2006).

In many cases, millimeter observations confirm the radial distribution of dust and gas inferred from the analysis of the disk spectral energy distribution (hereafter SED). The large majority of pre-main-sequence stars exhibit a flux excess over the stellar photosphere from near-infrared to centimeter wavelengths. This excess arises from a “classical,” or primordial, disk that extends from few stellar radii to few hundreds of astronomical units. However, about 10% of the observed systems (Muzerolle et al. 2010) exhibit a deficit of flux in the near- and mid-infrared compared to a “classical” disk, but have similar level of far-infrared emission (i.e., “transition” disk; Strom et al. 1989). The SED shape suggests that these stars lack the emission from warm dust close to the central star. The lack of warm dust emission may result from several different effects such as the presence of a stellar companion, the formation of a planetary system, disk photoevaporation, magnetorotational instabilities, disk viscosity, and grain growth (Alexander & Armitage 2007; Chiang & Murray-Clay 2007; Ireland & Kraus 2008; Isella et al. 2009; Strom et al. 1989; Tanaka et al. 2005). Subsequent high-resolution millimeter-wave images have confirmed that several “transition” disks are indeed dust-depleted within the radius predicted by the SED modeling (Piétu et al. 2006; Brown et al. 2009; Hughes et al. 2009).

The increasing sensitivity and resolution of (sub)millimeter-wave interferometers are now providing unexpected insights about the structure of “classical” disks. Subarcsecond observations of AB Aur and RY Tau systems indeed show that these two circumstellar disks may be dust depleted within orbital radii of 70 AU and 13 AU, respectively (Piétu et al. 2005; Isella et al. 2010), although no disk clearing has been inferred from the analysis of the SED (Robitaille et al. 2007; Dullemont et al. 2001). It is therefore possible that other “classical” disks may be characterized by a disk surface density that drastically deviates from that inferred from the SED modeling. In particular, since the dust opacity in the infrared is much larger than at millimeter-wavelengths, it is reasonable to assume that a large fraction of “classical” disks might be characterized by partially cleared inner cavities which are not detectable through the analysis of the infrared SED.

In this paper, we discuss the Herbig Ae star MWC 758 (HD 36112), a 3 Myr old pre-main-sequence A5 star located near the edge of the Taurus star-forming region at a distance of 200$^{+50}_{-40}$ pc (Perryman et al. 1997; van den Ancker et al. 1998). Previous investigations based on the SED modeling and on spatially resolved interferometric observations suggested that...
than the 12CO (2–1) emission (Chapillon et al. 2008). This shows the orientation of the disk major axis derived from the CO observations after correcting for the proper motion as described in Section 3. The dashed line into two peaks in the continuum emission, strongly suggesting that the disk surface density decreases within an orbital radius of about 100 AU. Similar to the cases of AB Aur and RY Tau, the MWC 758 inner disk may be dust-depleted with respect to a power-law surface density profile, though no disk clearing has been inferred from the analysis of the SED.

The observations, presented in Sections 2 and 3, are compared with theoretical disk models to constrain the surface density profile of MWC 758 disk. The results are presented in Section 4. The discussion and conclusions follow in Section 5.

2. OBSERVATIONS AND DATA REDUCTION

2.1. SMA Observations

The MWC 758 system was observed at 345 GHz with the Submillimeter Array7 (SMA) on Mauna Kea, Hawaii, on 2008 January 28. The array was in the extended configuration, and the seven available antennas provided baseline lengths from 31 to 225 m. The atmospheric transparency was excellent through the track, with an opacity of 0.04 at 225 GHz measured at the nearby Caltech Submillimeter Observatory. The correlator provided 2 GHz of bandwidth in two sidebands separated by 10 GHz, with the upper sideband centered near the CO J = 3–2 line at 345.796 GHz. This line was observed at a spectral resolution 0.18 km s\(^{-1}\) in a single correlator “chunk” spanning 104 MHz. The primary flux calibrator was Titan, and the passband calibrators were the quasars 3C273 and 3C111. Complex gain calibration was carried out with observations of the quasar J0530+135 (1.75 Jy) interleaved with MWC 758. Data were obtained for MWC 758 over the hour angle range −3.2 to +4.7. All of the data editing and calibration were performed using standard routines in the MIR software package for IDL. The resulting map of the dust continuum emission is shown in Figure 1. Applying natural weighting to the data, we achieve a noise level of 2.1 mJy beam\(^{-1}\) and synthesized beam FWHM of 0.′76 × 0.′56. Figure 2 shows the CO emission in the velocity range between 2.94 and 8.93 km s\(^{-1}\), as obtained by averaging the observations every two channels. The resulting noise level is 83 mJy beam\(^{-1}\).

2.2. CARMA Observations

The dust thermal emission toward MWC 758 was observed on 2009 January using the Combined Array for Research in Millimeter-wave Astronomy8 (CARMA) located at the altitude of 7200 feet in the Inyo mountains of eastern California. The array was in the extended B configuration, which provides baselines between 80.9 and 946.3 m. The observations were obtained at the local oscillator frequency of 97.5 GHz (\(\lambda \sim 3.3 \text{ mm}\)). The CARMA correlator was configured with six bands, each of which was configured to 468 MHz bandwidth to provide maximum continuum sensitivity. No molecular emission line was observed with CARMA during this track. The passband shape was calibrated by observing 3C454.3 and the flux calibration was set by observing Uranus. The

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Note that Chapillon et al. and Eisner et al. assume a stellar distance of 140 and 150 pc, respectively. The radii quoted from these two papers have been therefore rescaled for the distance of 200 pc.

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Figure 2. Channel maps for the $^{12}$CO $J = 3–2$ line emission observed toward MWC 758 in the velocity range between 2.94 and 8.93 km s$^{-1}$. The contours are spaced by $3\sigma = 0.25$ Jy beam$^{-1}$. The figure also shows the residuals obtained by subtracting the best-fitting model CO emission from the observations. The symbol “⋆” marks the position of the central star after correcting for the proper motion.

(A color version of this figure is available in the online journal.)

Figure 3. Left panel: $^{12}$CO $J = 3–2$ velocity gradient across MWC 758 circumstellar disk. Red and blue colors indicate the redshifted and blueshifted disk velocities, respectively. Velocity contours start at 4.7 km s$^{-1}$ and are spaced by 0.3 km s$^{-1}$. The dashed line indicates the position angle of the apparent disk major axis, which is rotated by 65$^\circ$ to east from north. Central panel: $^{12}$CO $J = 3–2$ integrated intensity map. Right panel: $^{12}$CO $J = 3–2$ spectrum obtained integrating over the region shown in the intensity map.

(A color version of this figure is available in the online journal.)

quasar J0510+180 was observed every 15 minutes to correct for atmospheric and instrumental effects. The data reduction and the image reconstruction were obtained using the MIRIAD software package. Figure 1 shows the resulting map, which is characterized by an rms noise of 0.25 mJy beam$^{-1}$ and a synthesized beam FWHM of 0.80 $\times$ 0.62 when natural weighting is applied to the data.

3. RESULTS

Figures 1–3 show the dust and CO emission toward MWC 758 as observed by pointing the telescope at the International Celestial Reference System coordinates of the star at the epoch J2000, R.A. = 05$^h$30$^m$27.530 ($\pm$12.89 mas) and decl. = +25$^\circ$19$'$.5708 ($\pm$5.57 mas). The position of the star at the date of the observation is shown by the symbol “⋆” and is calculated by applying a proper motion of $\delta$(R.A.) = 5.23 $\pm$ 1.40 mas yr$^{-1}$ and $\delta$(decl.) = $-25.95 \pm 0.66$ mas yr$^{-1}$ to the J2000 coordinates (both the stellar coordinates and the proper motion are from Perryman et al. 1997). The uncertainties on the stellar position resulting both from the Hipparcos astrometry and proper motion uncertainties are smaller than the size of symbol itself.

The dust emission at both 0.87 and 3.3 mm is characterized by two peaks which are symmetric with respect to the position of the star and are separated by about 0.7, or 140 AU at the distance of 200 pc. At the wavelength of 0.87 mm, the northern and southern peaks have fluxes of 89 and 77 mJy beam$^{-1}$.
respectively. A similar flux ratio between the two peaks is measured at 3.3 mm, suggesting that the observed morphology corresponds to an asymmetric dust emission between the two sides of the disk. By integrating over a circle with a radius of 1.5 centered at the position of the star, we measure fluxes of 217 ± 40 mJy and 4.9 ± 1.7 mJy at 0.87 and 3.3 mm, respectively, where the errors account for both the thermal noise and absolute calibration uncertainties. The integrated fluxes are in agreement with the values measured at 2' angular resolution by Chapillon et al. (2008) suggesting that the dust emission is all coming from the spatial scales probed by our observations. Assuming that the disk millimeter emission follows a power law \( F(\nu) \propto \nu^{a} \), we derive a spectral index \( a_{\nu=0.87-3\,\text{mm}} = 3.0 \).

Figure 2 shows the \(^{12}\text{CO} J = 3–2\) channel maps in the velocity range between 2.94 and 8.93 km s\(^{-1}\). The CO emission is spatially resolved and shows a velocity pattern typical of an inclined rotating disk. The first moment of the CO emission, i.e., the velocity gradients across the disk, indicates that the western and eastern sides of the disk are, respectively, redshifted and blueshifted with respect to the relative velocity of the star (see the left panel of Figure 3). The velocity gradient also suggests a position angle of the disk major axis of about 65° as measured from north to east. The central panel of Figure 3 shows the spectrally integrated CO emission. In contrast to what is observed in the dust continuum, the CO emission has a single peak coincident with the position of the star. The CO emission is observed up to a distance of 1.3 in the redshifted side of the disk and only to about 1'' on the other side. The line integrated flux is 12.9 Jy km s\(^{-1}\). This value is about 14% smaller than that measured in single dish observations (Dent et al. 2005), suggesting that part of the CO emission is resolved out in the interferometric observations. As a consequence, the observed asymmetry in the CO emission might be due to the lack of short baselines in the SMA observations instead of to a real asymmetry in the gas temperature or distribution. Finally, the spatially integrated CO spectrum (the right panel of Figure 3) has a double peak profile centered at the velocity of \( v_{\text{lsr}} = 6.0 \) km s\(^{-1}\). The line emission is observed in the velocity range \( v_{\text{lsr}} \pm 4 \) km s\(^{-1}\) with a brighter flux in the redshifted half of the disk.

To date, the different morphologies observed in the dust and gas emission toward MWC 758 represent a unique case among the sample of circumstellar disks spatially resolved with millimeter-wave interferometers. Our observations suggest that the optical depth of the dust continuum emission decreases within a radius of 70–100 AU from the star. Despite that, the disk remains gas-rich down to at least a distance of 30 AU to explain the centrally peaked CO emission. In addition, the SED of MWC 758 shows a flux excess above the stellar photosphere between near- and far-infrared wavelengths typical of optically thick flaring disk extending to few stellar radii (see Figure 4). The following section is devoted to reconciling all these observations in the framework of existing disk models in order to constrain the radial distribution of the circumstellar material around MWC 758.

4. DATA ANALYSIS

4.1. Disk Models

To analyze the data, we adopt the disk model discussed in Isella et al. (2007, 2009). In brief, we compute the radial dust temperature profile following the “two-layer” approximation of Chiang & Goldreich (1997), where the disk is characterized by a warm surface layer and a cooler interior. Both temperatures are calculated as a function of the orbital radius by iterating on the vertical disk structure (see Dulullemond et al. 2001). The disk is in hydrostatic equilibrium between the gas pressure and the stellar gravity, which leads to a flared geometry with the opening angle increasing with the distance from the central star.

The dust opacity is calculated by assuming an interstellar grain composition (Pollack et al. 1994) and a particle size distribution between a minimum and a maximum value \( a_{\text{min}} \) and \( a_{\text{max}} \), according to \( n(a) \propto a^{-q} \). We fix \( a_{\text{min}} = 0.05 \mu\text{m} \) and vary \( a_{\text{max}} \) and \( q \) to reproduce the spectral index \( \alpha \) of the millimeter disk emission. For the sake of simplicity, we assume that the dust opacity is constant throughout the disk (although see Birnstiel et al. 2010).

The radial distribution of the circumstellar material follows the similarity solution for the disk surface density of a viscous Keplerian disk (Lynden-Bell & Pringle 1974) expressed by:

\[
\Sigma(r, t) = \Sigma_{0} \left( \frac{r}{r_{0}} \right)^{\gamma} \times \exp \left\{ -\frac{1}{2(2-\gamma)} \left[ \left( \frac{r}{r_{1}} \right)^{2(2-\gamma)} - 1 \right] \right\},
\]

where the characteristic radius \( r_{0}, \gamma, \) and the surface density normalization \( \Sigma_{0} \) are free parameters of the model, as well as the disk inclination and position angle.

From the derived dust density, temperature, and opacity, we calculate the disk SED and synthetic disk images in the dust continuum at 0.87 mm and 3.3 mm using the radiative transfer solution discussed in Dulullemond et al. (2001). The synthetic disk images are then Fourier transformed and sampled at the appropriate positions on the \((u, v)\) plane corresponding to our CARMA and SMA observations.

The CO emission is calculated assuming that rotational levels are thermalized. The CO temperature is parameterized as a function of the radius as \( T_{\text{CO}}(r) = T_{\text{CO}}(r_{0})(r/r_{0})^{\xi} \), where \( T_{\text{CO}}(r_{0}) \) and \( \xi \) are free parameters and \( r_{0} = 50 \) AU. The CO density is then calculated by assuming that the dust and the gas are well mixed with a gas-to-dust ratio of 100. The adopted CO abundance is \(^{13}\text{CO}/\text{H}_{2} = 7.0 \times 10^{-5}\). Finally, the CO channel maps are calculated in the velocity range probed by our observations by assuming that the emission originates...
from an inclined Keplerian disk. Both the disk inclination and the mass of the central star are free parameters of the model (see the Appendix of Isella et al. 2007, for more details on the CO emission model). As for the case of the dust emission, the synthetic CO channel maps are then Fourier transformed and compared to the measured complex visibilities.

The best-fitting models for both the dust and gas emission, and the uncertainties on the model free parameters, are then obtained by \( \chi^2 \) minimization carried out using the Markov Chain Monte Carlo method discussed in Isella et al. (2009).

4.2. Results of the Model Fitting

A satisfactory fit of the spatially resolved dust and CO emission, as well as of the SED between far-infrared and (sub)millimeter wavelengths, is obtained with the disk model presented in Table 1. The disk surface density is mainly constrained by the optically thin dust emission at 0.87 and 3.3 mm and is characterized by \( \gamma = -4.1 \) and \( r_t = 69 \) AU. As shown in Figure 4, \( \Sigma(r) \) has a “bell” shape that reaches the maximum value of 3.3 g cm\(^{-2}\) at the radius of 100 AU. For \( r < r_t \), the surface density is proportional to \( r^4 \), suggesting that a steep radial density profile is required to explain the two peaks observed in the dust emission. For \( r > 100 \) AU, \( \Sigma(r) \) fades exponentially as prescribed by Equation (1). Given the angular resolution and sensitivity of our observations, we can constrain the surface density only between 40 and 130 AU. In this range, the disk is optically thin to the dust continuum emission at both 0.87 and 3.3 mm. Within 40 AU, the radial profile of \( \Sigma(r) \) might significantly deviate from that shown in Figure 4. However, the disk must remain optically thin to the dust continuum millimeter-wave emission down to at least few AU in order to reproduce the two peaks shown in Figure 1.

Radially integrating the surface density, we obtain a disk mass of 0.01 \( M_\odot \). This result is affected by a large error due to the uncertainties on the dust opacity and on the gas-to-dust ratio. In particular, the analysis of the \( ^{12}\text{CO} \) (2–1) and (1–0) line emission suggests that the gas-to-dust ratio in MWC 758 disk might be 0.1, leading therefore to a smaller disk mass.

Despite the steep surface density profile, the \( ^{12}\text{CO} \) (3–2) emission remains optically thick down to a radius of about 20 AU. As a consequence, the CO emission does not constrain the disk surface density but it provides a measure of the CO temperature in the form \( T_\text{CO}(r) \propto 53 \text{ K} \times (50\text{ AU}/r)^{-0.1} \). Note that the \( ^{12}\text{CO} \) (3–2) emission remains optically thick even if the CO is largely depleted with respect to the standard abundance. In particular, assuming a CO abundance of \( 10^{-6} \) (Chapillon et al. 2008) we calculate a mean optical depth of about \( 10^3 \). The almost radially constant profile of the CO temperature is in good agreement with the value \( \zeta = 0.05 \pm 0.2 \) derived from the analysis of the \( ^{12}\text{CO} \) (2–1) line emission by Chapillon et al. (2008). In the assumption that the disk is in Keplerian rotation, we derive a disk inclination of 22° and a dynamical mass for the central star of 2.0 \( M_\odot \), similar to the value of 2.3 \( M_\odot \) derived from the H-R diagram using stellar evolutionary tracks (Manoj et al. 2006).

The residuals obtained by subtracting the best-fitting model to the observations reveal asymmetries in the MWC 758 disk. In the case of the dust emission (see Figure 1), our disk model fails to reproduce the flux measured in the northern side of the disk at both 0.87 and 3.3 mm. From the intensity of the residuals, we calculate that the dust density, opacity, and/or temperature in the northern side of the disk might be 20% higher than the southern side. Figure 2 shows the residuals for the CO channel maps in the velocity range between 2.94 and 6.11 km s\(^{-1}\). Most of the significant structures appear in low-velocity channels of the redshifted side of the disk between 6.11 and 7.11 km s\(^{-1}\). Since the CO emission is almost completely optically thick, this suggests that the CO temperature might be higher in this side or that the CO extends to slightly larger distances from the central star. Note, however, that the CO asymmetry might also be due to the lack of short spatial frequency in SMA observations as discussed in Section 3.

Figure 5 shows the comparison between the observed SED and the SED of the best-fitting model to the millimeter data. At millimeter wavelengths, a good agreement with the observations is obtained with a grain size distribution slope \( q = 3.5 \) and a maximum grain size \( a_{\text{max}} = 1 \) cm. Fitting the dust opacity \( k(\lambda) \) between 1.3 and 2.8 mm with a power law \( k(\lambda) \propto \lambda^{-\beta} \), we obtain \( \beta = 1.35 \), which is a factor of two higher than the mean value measured in T Tauri stars but in agreement with those found in other Herbig Ae stars (see, e.g., Ricci et al. 2010; Natta et al. 2004).

The good agreement between the model and the observation extends down to a wavelengths of about 20 \( \mu \text{m} \). This suggests that the disk has a flared structure in the region that emits most of the mid- and far-infrared emission, namely between 20 and 110 AU. At shorter wavelengths, the emission of the best-fitting model is much lower than what observed toward MWC 758. This is due to the fact that for \( \Sigma(r) \propto r^4 \), the dust density within few AU from the central star is lower than \( 10^{-6} \text{ g cm}^{-2} \) and its near-infrared emission is negligible with respect to the stellar photospheric emission. The observed near-

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**Table 1**

Best-fitting Model Parameters for the Dust and CO Millimeter Emission Observed Toward MWC 758

| Disk orientation: | Inclination = 21° ± 2°, P.A. = 65° ± 7° |
|-------------------|----------------------------------------|
| Surface density \( \Sigma(r) \): | \( \gamma = -4.1 \pm 0.3, r_t = 69 \pm 3 \text{ AU}, \Sigma_t = 1.5 \pm 0.1 \text{ g cm}^{-2} \) |
| CO temperature \( T(\text{CO}) \): | \( T_{\text{CO}}(50 \text{ AU}) = 53 \pm 4 \text{ K}, \zeta = -0.1 \pm 0.2 \) |
| \( M_\star \): | \( 2.0 \pm 0.2 \text{ } M_\odot \) |
| Grain size distribution \( n(a) \): | \( a_{\text{max}} = 1 \text{ cm}, q = 3.5 \) |

**Note.** The quoted uncertainties correspond to the 1σ uncertainty level.
infrared excess therefore suggests the presence of material with temperature between 1000 and 2000 K in excess to what was predicted by our model.

To reconcile the near-infrared SED with our millimeter-wave observations, we propose two different scenarios. In the first case, the near-infrared emission is coming from an optically thick disk extending between 0.05 and 5 AU. This model comprises a gaseous disk between 0.05 and 0.1 AU, a “puffed-up” inner rim of the dusty disk between 0.4 and 0.5 AU (see, e.g., Isella & Natta 2005), and a flaring disk between 0.5 and 5 AU. As discussed in Isella et al. (2008), this model explains the spectro-interferometric observations of the H- (1.6 μm) and K-band (2.2 μm) disk emission (see also Benisty et al. 2010). The resulting SED, which is indicated by the short-dashed green line in Figure 5, provides a good description of the near-infrared emission and of the 10 μm silicate feature. Alternatively, the infrared SED can be explained with the presence of an optically thin layer of dust surrounding the star (Vinković et al. 2006; Mulders et al. 2010). In this case, we calculate that a ring of dust extending between 0.15 and 0.8 AU and characterized by an optical depth for the stellar radiation along the direction perpendicular to the disk mid-plane of 0.4 is required to reproduce the disk emission (see the long-dashed curve in Figure 5). For the assumed dust opacity, this model requires a minimum surface density of 0.01 g cm$^{-2}$ or a total mass of $2 \times 10^{-5} M_\odot$. Note that neither of these models will produce a significant dust and CO emission at millimeter wavelengths.

To summarize, assuming the similarity solution for a Keplerian viscous disk, we find strong evidence that the surface density of the MWC 758 disk increases as $r^q$ between 40 and 100 AU from the central star and decreases exponentially outward. However, this solution for the surface density is not unique and observations with higher angular resolution are required to discriminate between different models. We note, for example, that a discontinuity in the surface density at about 80 AU from the central star might provide an equally good fit with a different step gradient. Within 40 AU, $\Sigma(r)$ is constrained by the analysis of the SED. We calculate that a minimum surface density value of $10^{-2}$ g cm$^{-2}$ between 0.15 and 0.8 AU is required to reproduce the observed near- and mid-infrared emission. Probing these spatial scales requires high sensitivity 0.1 resolution observations that will become possible with the Atacama Large Millimeter Array (ALMA).

5. DISCUSSION AND CONCLUSIONS

The surface density profile in MWC 758 is similar to that measured in the transitional disk around the T Tauri star LkCa 15. In this case, $\Sigma(r)$ peaks at about 50 AU from the central star (Piétu et al. 2006; Espaillat et al. 2007) and a minimum mass of $1 \times 10^{-11} M_\odot$ within 5 AU from the central star is required to produce the observed near-infrared emission (Mulders et al. 2010). Similarly to what suggested for LkCa15, the shape of the surface density in MWC 758 may be due to the disk clearing operated by a low-mass stellar companion or by one or more giant planets. To this regard, we note that recent spectroastrometric observations of MWC 758 exclude the presence of a stellar companion more massive than about 0.9 $M_\odot$ orbiting between 0.15 and 0.3, or 30 and 60 AU for the assumed stellar distance of 200 pc (Wheelwright et al. 2010; H. E. Wheelwright 2010, private communication). The nature of a possible companion of MWC 758 remains therefore uncertain.

The asymmetries observed in the dust and CO emission support the hypothesis that the disk may be gravitationally perturbed by a massive body. The asymmetric CO emission can be due to a warped optically thick inner disk, i.e., more inclined than the outer disk, so that the CO in northwest side of the disk is partially shielded by the inner disk from the stellar radiation. This would result in a lower gas temperature and consequently lower emission with respect to the other side of the disk (see, e.g., Panić et al. 2010). Indeed, we note that the inclination of the inner disk derived from near-infrared interferometric observations is of about $40^\circ \pm 10^\circ$ (Isella et al. 2006, 2008; Eisner et al. 2004), which is larger than the $21^\circ \pm 2^\circ$ derived in this work.

Other mechanisms, which are in principle able to evacuate the inner disk, do not seem to be consistent with the observations. Photoevaporation and magnetorotational instability models predict completely cleared inner disks and surface density profiles at the inner disk radius much steeper than $r^q$ (Alexander & Armitage 2007; Chiang & Murray-Clay 2007). In addition, the surface density within few AU is much larger than that predicted by the viscous solution for the surface density that best fits the millimeter data. This fact excludes that $\Sigma(r)$ is shaped only by the radial profile of the disk viscosity (see the discussion in Isella et al. 2009). Low dust opacity caused by grain growth has also been suggested to explain inner holes observed in the continuum emission at millimeter wavelengths (Tanaka et al. 2005). In the case of MWC 758, the dust opacity should decrease by at least two orders of magnitude to explain the observations. This would require a grain size distribution flatter than what typically assumed (i.e., $n(a) \propto a^{-q}$ with $q < 3$) and a steep radial decrease of the maximum grain (i.e., from 10 cm at 10 AU to 10 μm at 100 AU).

In conclusion, our observations show that the disk around MWC 758 is characterized by a steeply increasing surface density between 40 and 100 AU from the central star. Given the infrared SED and the asymmetries observed in the gas and dust.
emission, we conclude that the surface density might be shaped by the presence of stellar or planetary size companion orbiting within 30 AU. The fact that MWC 758 was previously labeled as “classical” disk emphasizes that SEDs alone do not provide a complete picture of disk structure and that high-resolution millimeter-wave images are essential to reveal the structure of the cool disk mid-plane.

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