EIGENPROBLEM FOR JACOBI MATRICES: 
HYPERGEOMETRIC SERIES SOLUTION

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Abstract. We study the perturbative power-series expansions of the eigenvalues and eigenvectors of a general tridiagonal (Jacobi) matrix of dimension $d$. The (small) expansion parameters are the entries of the two diagonals of length $d - 1$ sandwiching the principal diagonal which gives the unperturbed spectrum.

The solution is found explicitly in terms of multivariable (Horn-type) hypergeometric series in $3d - 5$ variables in the generic case. To derive the result, we first rewrite the spectral problem for the Jacobi matrix as an equivalent system of algebraic equations which are then solved by the application of the multivariable Lagrange inversion formula. The corresponding Jacobi determinant is calculated explicitly. Explicit formulae are also found for any monomial composed of eigenvector’s components.

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1. Introduction

The problem of solving algebraic equations by series expansions has a long history. In case of a single polynomial equation a major cornerstone is Birkeland’s paper [1] where a solution is given through an application of the Lagrange inversion formula [2]. Later on, an alternative approach was suggested by Mayr in [3], who derived the relevant hypergeometric series as solutions of some PDE’s satisfied by the zeros as functions of the coefficients of the polynomial. A modern interpretation in terms of Gel’fand-Kapranov-Zelevinsky hypergeometric functions [4, 5, 6] can be found in [7]. The latter approach can also be applied to general systems of algebraic equations.

The main goal of this paper is to derive a complete power-series solution of the spectral problem for a finite $(d \times d)$ Jacobi matrix $M$. We consider the off-diagonal matrix elements to be small, so that the whole problem looks like a perturbation of a diagonal matrix. After fixing a normalization of the eigenvector the problem is reduced to solving a system of $d$ quadratic equations for $d$ unknowns. Then we transform it into an equivalent larger system of special (Lagrange-form) $3d - 5$ cubic equations, which are then inverted by the application of the multivariable Lagrange
functions of the parameters $\alpha, \beta, \gamma$ produces a system of algebraic (quadratic) equations for the eigenvalue $\Lambda_k$ and the components of the eigenvector $V_j^{(k)}$, respectively, where the components $V_j^{(k)}$ of $V^{(k)}$ can be chosen as $V_j^{(k)} = \delta_{jk}$. By continuity argument, for the small values of $\beta_k, \gamma_k$ the eigenvalues $\Lambda_k, k = 1, \ldots, d$, are distinct and can be numbered in such a way that

$$\Lambda_k = \alpha_k + \text{higher order terms.}$$

We also choose to normalise the eigenvector $V^{(k)}$ by the condition

$$MV^{(k)} = \Lambda_k V^{(k)}$$

for its $k^{th}$ component. Therefore, the remaining components must vanish in the zeroth approximation:

$$V_j^{(k)} = 0 + \text{higher order terms, \quad} j \neq k.$$  

The eigenvalue problem together with the normalization condition produces a system of $d$ algebraic (quadratic) equations for the eigenvalue $\Lambda_k$ and the free components $V_j^{(k)}, j \neq k$ of the eigenvector $V_j^{(k)}$, defining them as algebraic functions of the parameters $\alpha, \beta, \gamma$. The conditions and fix uniquely the branches of the multivalued algebraic functions for small values of $\beta$ and $\gamma$.

The problem we solve in the present paper is to find an effective way to construct explicit expressions for the coefficients of the power series expansions for the eigenvalues $\Lambda_k$ and the components of the eigenvectors $V_j^{(k)}$. The available literature on solving systems of algebraic equations by multivariate hypergeometric series (see a review in ) focuses mainly on solving the generic algebraic system. We are not aware of any detailed analysis of the particular systems arising from the Jacobi matrix spectral problem. The importance of the latter problem for numerous applications in mathematical physics and, in particular, in the theory of quantum integrable systems has been the main motivation of our study. Rather then using more modern approaches, in the present paper we follow the original idea of Birkeland and use a variant of the Lagrange inversion formula.

We shall use the following variant of the multivariable Lagrange inversion theorem. Let boldface letters denote vectors $\xi \equiv (\xi_1, \ldots, \xi_D)$, multi-indices $\mathbf{q} \equiv (q_1, \ldots, q_D) \in \mathbb{Z}_+^D$, and monomials $\xi^\mathbf{q} \equiv \xi_1^{q_1} \cdots \xi_D^{q_D}$. The inequality $\mathbf{q} \geq \mathbf{p}$ is understood component-wise: $q_i \geq p_i, \forall i$. Let $[\xi^\mathbf{q}]h(\xi)$ denote the coefficient at $\xi^\mathbf{q}$ in the power series $h(\xi)$.
Theorem 1.1. Let \( \xi = (\xi_1, \ldots, \xi_D) \), \( \eta = (\eta_1, \ldots, \eta_D) \). Let \( \varphi_i(\eta) \), \( i = 1, \ldots, D \), be formal power series in \( \eta \) such that \( \varphi_i(0) \neq 0 \) \( \forall i \). Then the system of \( D \) equations
\[
\xi_i = \frac{\eta_i}{\varphi_i(\eta)}, \quad i = 1, \ldots, D,
\]

defines uniquely \( \eta_i(\xi), i = 1, \ldots, D \), as formal power series in \( \xi \).

In addition, let \( \chi(\eta) \) be a multiple Laurent series, that is, \( \chi(\eta)\eta^p \) is a power series for some \( p \geq 0 \). Then the Laurent series expansion for \( \chi(\eta(\xi)) \)
\[
\chi(\eta(\xi)) = \sum_{q \geq -p} X_q \xi^q
\]
is given by the formula
\[
X_q \equiv [\xi^q] \chi(\eta(\xi)) = [\eta^q] \chi(\eta) \varphi^q(\eta) J(\eta),
\]
where \( J \) is the Jacobian
\[
J = \det \left(1 - \frac{\partial \log \varphi}{\partial \log \eta}\right) = \det \left| \delta_{jk} - \frac{\eta_k}{\varphi_j} \frac{\partial \varphi_j(\eta)}{\partial \eta_k} \right|.
\]

The analytic version of the Lagrange theorem guarantees that if \( \varphi_i(\eta) \) and \( \chi(\eta)\eta^p \) are analytic at 0 then the Laurent series converges in a punctured neighbourhood around \( \xi = 0 \).

When applying the Lagrange formula, the major complication comes from the Jacobian \( J \), which may be difficult to compute. Fortunately, for our particular problem the Jacobian can be calculated explicitly, in a relatively compact form.

The paper is organised as follows. In section 2 we write down the set of quadratic equations defining the eigenvalue \( \Lambda \) and the components \( V_j \) of the eigenvector \( V \), transform them into the form which is convenient for study, identify the combinations of the small parameters which serve as the expansion variables, and rewrite the equations again in the form which allows us to apply Lagrange’s inversion formula. In section 3 we calculate an important component of Lagrange’s formula: the Jacobian \( J \).

In section 4 we put together all the ingredients of the Lagrange formula and produce explicit expressions for all \( \Lambda_k \)’s and \( V_j^{(k)} \)’s as finite sums of power series. The number of terms in the sum equals to the number of terms in the Jacobian \( J \). All the power series are particular cases of a single universal Horn-type multivariable hypergeometric series which we denote \( \Phi \). In generic situation, the series \( \Phi \) depends on \( 3d - 5 \) expansion variables and \( d - 1 \) integer parameters. In section 5 we describe the simplification of our results for the special case \( k = 1 \) (or \( k = d \)) when the eigenvalue \( \Lambda \) stems from a corner of the matrix \( M \). In section 6 we examine a few low-dimensional examples illustrating the general results. The last section contains a discussion of possible applications and extensions of our result.

2. Lagrange Equations

From now on we shall concentrate on studying a single eigenvalue \( \Lambda_k \) and the corresponding eigenvector \( V_j^{(k)} \), for some fixed value of the index \( k \). We shall change our notation accordingly, to simplify the calculations. Set \( r = d - k \) and \( \bar{r} = k - 1 \), so that
\[
d = r + \bar{r} + 1.
\]

Let \( \Lambda_k = \alpha_k + \lambda \), \( V_j^{(k)} = \tilde{v}_{k-j} \) and \( \alpha_j = \tilde{a}_{k-j} \) for \( j = 1, \ldots, k \), and \( V_j^{(k)} = v_{j-k} \) and \( \alpha_j = a_{j-k} \) for \( j = k, \ldots, d \), so that \( V_j^{(k)} = v_0 = \tilde{v}_0 \) and \( \alpha_k = a_0 = \tilde{a}_0 \).

Respectively, let \( \beta_j = \tilde{b}_{k-j-1} \) and \( \alpha_j = \tilde{a}_{k-j-1} \) for \( j = 1, \ldots, k - 1 \), and \( \beta_j = \beta_{j-k} \) and \( \alpha_j = \alpha_{j-k} \) for \( j = k, \ldots, d - 1 \). Without loss of generality we can set \( \alpha_k = \ldots = \alpha_{k-1} = \ldots = \alpha_{d-1} = \ldots = \alpha_0 = 0 \).
$a_0 = \bar{a}_0 = 0$. The normalisation condition \[\text{(1.4)}\] implies that $V^{(k)}_k = v_0 = \bar{v}_0 = 1$.

Since $\alpha_i$ are assumed distinct, we have $a_i \neq 0$ and $\bar{a}_i \neq 0$ for $i \neq 0$.

As a result, the eigenvalue problem \[\text{(1.3)}\] takes the following form:

$$
\begin{pmatrix}
\tilde{a}_r & b_{r-1} \\
\cdots & \cdots \\
\tilde{c}_i & \tilde{a}_i & b_{i-1} \\
\cdots & \cdots \\
\tilde{c}_1 & \tilde{a}_1 & b_0 \\
0 & a_1 & b_1 \\
\cdots & \cdots \\
1 & a_i & b_i \\
\cdots & \cdots \\
1 & a_r & b_r \\
\end{pmatrix}
\begin{pmatrix}
v_r \\
\cdots \\
v_i \\
\cdots \\
v_1 \\
1 \\
\cdots \\
v_i \\
\cdots \\
v_r \\
\end{pmatrix}
= \lambda
\begin{pmatrix}
\tilde{v}_r \\
\cdots \\
\tilde{v}_i \\
\cdots \\
\tilde{v}_1 \\
1 \\
\cdots \\
\tilde{v}_i \\
\cdots \\
\tilde{v}_r \\
\end{pmatrix},
$$

(2.2)

The numbers $\bar{r}, r = 0, 1, \ldots$ measure the distances of the selected diagonal element from the corners of the matrix. The cases $\bar{r}, r = 0, 1$ are slightly special, we shall comment on them in due course, all other cases with $\bar{r}, r \geq 2$ are generic.

Expressing $\lambda$ from the ‘central’ (‘zeroth’) row as

$$
\lambda = \tilde{c}_0 \tilde{v}_1 + b_0 v_1
$$

(2.3)

and substituting it into the remaining rows we get the set of $r + \bar{r} = d - 1$, see \[\text{(2.1)}\], quadratic equations

$$
c_{i-1}v_{i-1} + a_i v_i + b_i v_{i+1} = \tilde{c}_0 \tilde{v}_1 v_i + b_0 v_1 v_i, \\
\tilde{c}_i \tilde{v}_{i+1} + \tilde{a}_i \tilde{v}_i + \tilde{b}_{i-1} \tilde{v}_{i-1} = \tilde{c}_0 \tilde{v}_1 \tilde{v}_i + b_0 v_1 \tilde{v}_i
$$

(2.4a)

(2.4b)

where we assume that $v_0 \equiv \tilde{v}_0 = 1, b_\bar{r} = 0, \tilde{c}_\bar{r} = 0$. Formula \[\text{(2.3)}\] eliminates the eigenvalue $\lambda$ by expressing it in terms of the two eigenvector’s components: $v_i$ and $\tilde{v}_i$. From now on, the variables $v_i$ and $\tilde{v}_i$ will be our only $d - 1$ unknowns.

Note that equations \[\text{(2.4)}\] are invariant with respect to the ‘involution’ (rotation by 180 degrees around the central element)

$$
r \leftrightarrow \bar{r}, \quad v_i \leftrightarrow \tilde{v}_i, \quad a_i \leftrightarrow \tilde{a}_i, \quad b_i \leftrightarrow \tilde{b}_i, \quad c_i \leftrightarrow \tilde{c}_i
$$

(2.5)

which we refer to as the tilde-symmetry.

The next step is to rescale the variables $v_i$ and $\tilde{v}_i$ in order to identify the convenient combinations of the expansion parameters. Setting $b_0 = \ldots = b_{\bar{r}-1} = 0$ and $\tilde{c}_0 = \ldots = \tilde{c}_{\bar{r}-1} = 0$ and denoting the corresponding values of $v_i$ and $\tilde{v}_i$ as, respectively, $v_i^0$ and $\tilde{v}_i^0$ we get the equations

$$
c_{i-1}v_{i-1}^0 + a_i v_i^0 = 0 \quad \Rightarrow \quad v_i^0 = -\frac{c_{i-1}}{a_i} v_{i-1}^0, \quad (2.6a)
$$

$$
\tilde{a}_i \tilde{v}_i^0 + \tilde{b}_{i-1} \tilde{v}_{i-1}^0 = 0 \quad \Rightarrow \quad \tilde{v}_i = -\frac{\tilde{b}_{i-1}}{\tilde{a}_i} \tilde{v}_{i-1}^0, \quad (2.6b)
$$

which can be solved recursively yielding

$$
v_i^0 = (-1)^i \frac{c_0 \ldots c_{i-1}}{a_1 \ldots a_i}, \quad i = 1, \ldots, r, \quad (2.7a)
$$

$$
\tilde{v}_i^0 = (-1)^i \frac{\tilde{b}_0 \ldots \tilde{b}_{i-1}}{\tilde{a}_1 \ldots \tilde{a}_i}, \quad i = 1, \ldots, \bar{r}. \quad (2.7b)
$$
Let us rescale \( v_i \) and \( \tilde{v}_i \) by the formulae
\[
 v_i = v_0^i u_i = (-1)^i \frac{c_0 \cdots c_{i-1}}{a_1 \cdots a_i} u_i, \quad i = 1, \ldots, r, \tag{2.8a}
\]
\[
 \tilde{v}_i = \tilde{v}_0^i \tilde{u}_i = (-1)^i \frac{\tilde{b}_0 \cdots \tilde{b}_{i-1}}{\tilde{a}_1 \cdots \tilde{a}_i} \tilde{u}_i, \quad i = 1, \ldots, \tilde{r}, \tag{2.8b}
\]
and rewrite equations (2.4) in terms of \( u_i, \tilde{u}_i \):
\[
u_i = u_{i-1} - \frac{b_0 c_0}{a_1 a_i} u_1 u_i - \frac{b_0 \tilde{c}_0}{a_1 a_i} \tilde{u}_1 \tilde{u}_i + \frac{b_i c_i}{a_i a_{i+1}} u_{i+1}, \quad i = 1, \ldots, r, \tag{2.9a}
\]
\[
\tilde{u}_i = \tilde{u}_{i-1} - \frac{\tilde{b}_0 c_0}{\tilde{a}_1 \tilde{a}_i} \tilde{u}_1 \tilde{u}_i - \frac{\tilde{b}_0 \tilde{c}_0}{\tilde{a}_1 \tilde{a}_i} \tilde{u}_1 \tilde{u}_i + \frac{\tilde{b}_i \tilde{c}_i}{\tilde{a}_i \tilde{a}_{i+1}} \tilde{u}_{i+1}, \quad i = 1, \ldots, \tilde{r}, \tag{2.9b}
\]
where we assume \( u_0 = \tilde{u}_0 = 1, b_r = \tilde{c}_r = 0 \).

Equations (2.8) contain the parameters \( a, b, c, \tilde{a}, \tilde{b}, \tilde{c} \) in specific combinations, which are convenient to use as the expansion parameters. Note that \( b_i \) and \( c_i \) enter only through the product \( b_i c_i \) (same for \( \tilde{b}_i \tilde{c}_i \)). Introduce the variables \( x_i, y_i, z_i \)
\[
x_i = \frac{b_0 c_0}{a_1 a_{i+1}}, \quad i = 0, \ldots, r - 1, \tag{2.10a}
\]
\[
y_i = \frac{b_0 \tilde{c}_0}{a_1 a_{i+1}}, \quad i = 0, \ldots, r - 1, \tag{2.10b}
\]
\[
z_i = \frac{b_{i+1} c_{i+1}}{a_{i+1} a_{i+2}}, \quad i = 0, \ldots, r - 2, \tag{2.10c}
\]
(altogether \( r + r + (r - 1) = 3r - 1 \) variables) as well as their tilde-analogs
\[
\bar{x}_i = \frac{\tilde{b}_0 \tilde{c}_0}{\tilde{a}_1 \tilde{a}_{i+1}}, \quad i = 0, \ldots, \tilde{r} - 1, \tag{2.11a}
\]
\[
\bar{y}_i = \frac{\tilde{b}_0 c_0}{\tilde{a}_1 \tilde{a}_{i+1}}, \quad i = 0, \ldots, \tilde{r} - 1, \tag{2.11b}
\]
\[
\bar{z}_i = \frac{\tilde{b}_{i+1} \tilde{c}_{i+1}}{\tilde{a}_{i+1} \tilde{a}_{i+2}}, \quad i = 0, \ldots, \tilde{r} - 2, \tag{2.11c}
\]
(altogether \( \tilde{r} + \tilde{r} + (\tilde{r} - 1) = 3\tilde{r} - 1 \) variables). The total number of variables is then \( D = 3r + 3\tilde{r} - 2 = 3d - 5 \).

Equations (2.4) are simplified now to the form
\[
u_i = u_{i-1} - x_{i-1} u_1 u_i - y_{i-1} u_1 u_i + z_{i-1} u_{i+1}, \quad i = 1, \ldots, r, \tag{2.12a}
\]
\[
\tilde{u}_i = \tilde{u}_{i-1} - \bar{x}_{i-1} \tilde{u}_1 \tilde{u}_i - \bar{y}_{i-1} \tilde{u}_1 \tilde{u}_i + \bar{z}_{i-1} \tilde{u}_{i+1}, \quad i = 1, \ldots, \tilde{r}, \tag{2.12b}
\]
where we assume \( u_0 = \tilde{u}_0 = 1, z_{r-1} = \tilde{z}_{\tilde{r}-1} = 0 \). Since \( b_i, \tilde{b}_i, c_i, \tilde{c}_i \) are small parameters, so are \( x_i, y_i, z_i, \bar{x}_i, \bar{y}_i, \bar{z}_i \). In the zeroth approximation we have a set of binomial equations \( u_i^0 = \tilde{u}_i^0, \bar{u}_i^0 = \bar{u}_i^{0} \). The iteration of equations (2.4) with the initial values \( u_i^0 = \bar{u}_i^0 = 1 \) produces formal power series expansions of \( u_i \) and \( \tilde{u}_i \) in the variables \( (x, y, z, \bar{x}, \bar{y}, \bar{z}) = \xi \).

The tilde-symmetry (2.84) for equations (2.2) takes the form
\[
r \leftrightarrow \tilde{r}, \quad u_i \leftrightarrow \tilde{u}_i, \quad x_i \leftrightarrow \bar{x}_i, \quad y_i \leftrightarrow \bar{y}_i, \quad z_i \leftrightarrow \bar{z}_i. \tag{2.13}
\]

**Important remark:** the variables \( \xi \) are, generally speaking, not independent. Indeed, the \( 3d - 5 \) variables are bound by the relations
\[
\frac{x_i}{y_i} = \frac{\bar{y}_i}{\bar{x}_j} = \frac{b_0 c_0 \bar{a}_1}{\tilde{b}_0 \tilde{c}_0 \tilde{a}_1}, \quad i = 0, \ldots, r - 1, \quad j = 0, \ldots, \tilde{r} - 1, \tag{2.14}
\]
altogether \( r + \tilde{r} - 1 = d - 2 \) independent relations, which leaves only \( 2d - 3 \) independent variables. Only in the special case \( r = 0 \) (or \( \tilde{r} = 0 \)), see end of this section, the
variables $\xi$ become independent. Nevertheless, when solving equations (2.12) by power series, it is convenient to treat $\xi$ as the set of $3d - 5$ independent variables. The expressions (2.10), (2.11) for $\xi$ in terms of $(a, b, c, \tilde{a}, \tilde{b}, \tilde{c})$ can be substituted then into the resulting expansions at the very final stage.

In order to find explicitly the coefficients of the power series for $u_i$ and $\tilde{u}_i$, we shall use the Lagrange inversion formula. In the notation of Theorem 1.1, the expansion

$$\eta^a$$

respectively, of the 6 matching sets of expandable quantities $\eta = (s, \tilde{s}, w, \tilde{w}, t, \tilde{t})$:

$$s_i \sim x_i, \quad i = 0, \ldots, r - 1; \quad \tilde{s}_j \sim \tilde{x}_j, \quad j = 0, \ldots, \tilde{r} - 1,$$

(2.15a)

$$t_i \sim y_i, \quad i = 0, \ldots, r - 1; \quad \tilde{t}_j \sim \tilde{y}_j, \quad j = 0, \ldots, \tilde{r} - 1,$$

(2.15b)

$$w_i \sim z_i, \quad i = 0, \ldots, r - 2; \quad \tilde{w}_i \sim \tilde{z}_i, \quad i = 0, \ldots, \tilde{r} - 2.$$ 

(2.15c)

Note that the variables $(u, \tilde{u})$ cannot be used directly as $\eta$ because they have nonzero limits for $\xi \to 0$, therefore they have to be completed by small factors. Besides, the number of the variables $(u, \tilde{u})$ is only $r + \tilde{r} = d - 1$, so we need $2d - 4$ extra variables to match the number $D = 3d - 5$ of expansion variables. Actually, it is more convenient to split the number $D$ differently: $D \equiv 3d - 5 = (2d - 2) + (d - 3)$, with the first set of $2d - 2$ variables for the larger system coming from only two variables, $u_1$ and $\tilde{u}_1$. Define $s_i, t_i, \tilde{s}_i, \tilde{t}_i$ as

$$s_i = x_i u_1, \quad t_i = \tilde{y}_i \tilde{u}_1, \quad i = 0, \ldots, r - 1,$$

(2.16a)

$$\tilde{s}_i = \tilde{x}_i \tilde{u}_1, \quad \tilde{t}_i = \tilde{y}_i u_1, \quad i = 0, \ldots, \tilde{r} - 1.$$ 

(2.16b)

The second set of $d - 3$ remaining variables $w, \tilde{w}$ is defined as

$$w_i = z_i \frac{u_i + 2}{u_i}, \quad i = 0, \ldots, r - 2,$$

(2.17a)

$$\tilde{w}_i = \tilde{z}_i \frac{\tilde{u}_i + 2}{\tilde{u}_i}, \quad i = 0, \ldots, \tilde{r} - 2.$$ 

(2.17b)

It might seem easier to define $w_i$ to be proportional to $u_i$ rather than to the ratio of two $u$’s with the step 2. However, a little experimenting shows that our choice of $w_i$ leads to simpler, factorised expressions for the functions $\varphi_i(\eta)$ in (2.16).

From $u_0 = \tilde{u}_0 = 1$ it follows that

$$w_0 = z_0 u_2,$$

(2.18a)

$$\tilde{w}_0 = \tilde{z}_0 \tilde{u}_2.$$ 

(2.18b)

Solving equations (2.17) – (2.19) recursively we obtain the expressions for $u_2, \ldots, u_r$ in terms of $w_j$ and for $\tilde{u}_2, \ldots, \tilde{u}_\tilde{r}$ in terms of $\tilde{w}_j$: 

$$u_{2j} = \frac{u_{2j} - 2 \ldots u_0}{z_{2j} - 2 \ldots z_0}, \quad u_{2j+1} = \frac{w_{2j-1} \ldots w_0 s_0}{z_{2j-1} - 2 \ldots z_1 x_0},$$

(2.19a)

$$\tilde{u}_{2j} = \frac{\tilde{w}_{2j} - 2 \ldots \tilde{u}_0}{z_{2j} - 2 \ldots z_0}, \quad \tilde{u}_{2j+1} = \frac{\tilde{w}_{2j-1} \ldots \tilde{w}_1 t_0}{z_{2j-1} - 2 \ldots z_1 \tilde{x}_0}.$$ 

(2.19b)
Proposition 2.1. The set of equations (2.12a), (2.16a) and (2.17a) is equivalent to the set of Lagrange-type (1.6) equations

\[ x_i = \frac{s_i}{f_i}, \quad y_i = \frac{t_i}{g_i}, \quad i = 0, \ldots, r - 1, \quad z_j = \frac{w_j}{h_j}, \quad j = 0, \ldots, r - 2, \]

(2.20a)

\[ \tilde{x}_i = \frac{\tilde{s}_i}{f_i}, \quad \tilde{y}_i = \frac{\tilde{t}_i}{g_i}, \quad i = 0, \ldots, \tilde{r} - 1, \quad \tilde{z}_j = \frac{\tilde{w}_j}{\tilde{h}_j}, \quad j = 0, \ldots, \tilde{r} - 2, \]

(2.20b)

where

\[ f_0 = \ldots = f_{r-1} = \tilde{g}_0 = \ldots = \tilde{g}_{\tilde{r}-1} = \frac{1 + w_0}{1 + s_0 + t_0}, \]  

(2.21a)

\[ h_i = \frac{(1 + w_i)(1 + w_{i+1})}{(1 + s_i + t_i)(1 + s_{i+1} + t_{i+1})}, \quad i = 0, \ldots, r - 2 \quad (w_{r-1} \equiv 0), \]  

(2.21b)

\[ \tilde{f}_0 = \ldots = \tilde{f}_{\tilde{r}-1} = g_0 = \ldots = g_{r-1} = \frac{1 + \tilde{w}_0}{1 + \tilde{s}_0 + \tilde{t}_0}, \]  

(2.22a)

\[ \tilde{h}_i = \frac{(1 + \tilde{w}_i)(1 + \tilde{w}_{i+1})}{(1 + \tilde{s}_i + \tilde{t}_i)(1 + \tilde{s}_{i+1} + \tilde{t}_{i+1})}, \quad i = 0, \ldots, \tilde{r} - 2 \quad (\tilde{w}_{\tilde{r}-1} \equiv 0). \]  

(2.22b)

In the notation of Theorem 1.1, we have

\[ \xi = (x, y, z, \tilde{x}, \tilde{y}, \tilde{z}), \quad \eta = (s, t, w, \tilde{s}, \tilde{t}, \tilde{w}), \quad \varphi = (f, g, h, \tilde{f}, \tilde{g}, \tilde{h}). \]  

(2.23)

Proof. Consider equation (2.12a) for \( u_{i+1} \):

\[ u_{i+1} = u_i - x_i u_1 u_{i+1} - y_i \tilde{u}_1 u_{i+1} + z_i u_{i+2}. \]  

(2.24)

Note that \( z_i u_{i+2} = w_i u_i \) by virtue of (2.17a). Having rearranged the terms we get

\[ u_{i+1}(1 + x_i u_1 + y_i \tilde{u}_1) = u_i(1 + w_i). \]  

(2.25)

From (2.16a) it follows that \( x_i u_1 = s_i \) and \( y_i \tilde{u}_1 = t_i \). Therefore,

\[ \frac{u_i}{u_{i+1}} = \frac{1 + s_i + t_i}{1 + w_i}. \]  

(2.26)

Multiplying the equalities (2.26) for the index \( i \) and \( i + 1 \) we get

\[ \frac{u_i}{u_{i+2}} = \frac{(1 + s_i + t_i)(1 + s_{i+1} + t_{i+1})}{(1 + w_i)(1 + w_{i+1})}. \]  

(2.27)

It remains to replace \( u_i/u_{i+2} \) with \( z_i/w_i \) from the equality (2.17a), and we obtain the equality of the form \( z_i = w_i/h_i \), see (2.20a), where

\[ h_i = \frac{(1 + w_i)(1 + w_{i+1})}{(1 + s_i + t_i)(1 + s_{i+1} + t_{i+1})}, \quad i = 0, \ldots, r - 2, \]  

(2.28)

and it is assumed that \( w_{r-1} = 0 \).

The special case of equation (2.12a) for \( i = 1 \),

\[ u_1 = 1 - x_0 w_1^2 - y_0 \tilde{u}_1 u_1 + z_0 u_2, \]  

(2.29)

has to be treated separately. Similarly to the general case, we replace \( z_0 u_2 \) with \( w_0 \) using (2.15a), and substitute \( u_1 = s_0/x_0 \) from (2.16a), and \( \tilde{u}_1 = \tilde{s}_0/\tilde{x}_0 \) from (2.16a). Then we replace \( y_0 \tilde{s}_0 \) with \( \tilde{x}_0 t_0 \), using both (2.16a) and (2.16e). As a result, the equation takes the form

\[ s_0(1 + s_0 + t_0) = x_0(1 + w_0). \]  

(2.30)
or, equivalently, the form \( x_0 = s_0/f_0 \), see (2.10a) for \( i = 0 \), where

\[
f_0 = \frac{1 + w_0}{1 + s_0 + t_0}.
\]

(2.31)

From (2.10a) and (2.10b) it also follows that

\[
f_i = \frac{s_i}{x_i} = u_i = \frac{t_i}{y_i} = g_i = \frac{s_0}{x_0} = f_0.
\]

(2.32)

The remaining half of the equations is obtained by the tilde-symmetry (2.13).

As was said before, the cases of a corner eigenvalue, i.e. when \( r = 0 \) or \( \tilde{r} = 0 \), and the next one, i.e. when \( r = 1 \) or \( \tilde{r} = 1 \), are slightly special. When one applies the formulae from the above theorem describing the generic case, i.e. when both \( r, \tilde{r} \geq 2 \), to such cases one has to remember that

- for a corner eigenvalue, say \( \tilde{r} = 0 \), there are no tilded variables and there are also no variables \( y, t, g \), so that in this case we only have \( 2d - 3 = 2(r + 1) - 3 = 2r - 1 \) variables, therefore one must disregard entirely and set \( t_j \equiv 0 \), \( j = 0, \ldots, r - 1 \), in (2.21);
- for the immediate next eigenvalue, say \( \tilde{r} = 1 \), all variables are present but one, \( w \), so that one must set \( \tilde{w}_0 \equiv 0 \) in (2.22a) and disregard (2.22b) entirely, there are \( 3d - 5 = 3(r + 2) - 5 = 3r + 1 \) variables in this case.

All these modifications are easily seen from the expressions (2.10d)–(2.11) of the small parameters \((x, y, z, \bar{x}, \bar{y}, \bar{z})\) in terms of the initial small parameters \((b, c, \hat{b}, \hat{c})\). For a more detailed study of the special cases see section 5.

3. Jacobian

In this section only, we use the ordering of variables which is different from the one used above: \( \xi = (\bar{x}, \bar{y}, x, y, z, \bar{z}) \), \( \eta = (s, t, s, t, \bar{w}, w) \), \( \varphi = (f, g, f, g, \bar{h}, h) \).

**Theorem 3.1.** The Jacobian \( J \) defined by (2.6) can be expressed as

\[
J = SS^\top - TT^\top,
\]

(3.1)

where

\[
S = \frac{1}{W} \left( 1 + \sum_{j=0}^{r-1} \frac{s_j(1 + w_j)}{1 + s_j + t_j} \prod_{k=0}^{j-1} w_k \right), \quad T = \frac{1}{W} \sum_{j=0}^{r-1} \frac{t_j(1 + w_j)}{1 + s_j + t_j} \prod_{k=0}^{j-1} w_k,
\]

(3.2)

and

\[
W = \prod_{k=0}^{r-2} (1 + w_k).
\]

(3.3)

Here we adopt the following agreement: whenever a sum has the upper limit smaller than the lower one then its value should be taken as 0. Also, any product in a similar case should be taken as 1. The tildes in \( S \) refer to replacing \( s_j, t_j, w_j \), and \( r \) by their tilded versions. As always, we assume that \( \tilde{w}_{r-1} \equiv 0 \) and \( w_{r-1} \equiv 0 \) (see Proposition 2.1).

Note that the number of terms in \( J \) is \( r\bar{r} + (r + 1)(\bar{r} + 1) = 2r\bar{r} + r + \bar{r} + 1 = 2dk - 2k^2 + 2k - d \) (recall that \( d = \bar{r} + r + 1 \) and \( k = \bar{r} + 1 \).

**Proof.** Consider the rows of the matrix \( J = \delta_{jk} - \frac{\partial f_j}{\partial x_k} \frac{\partial g_k}{\partial x_j} \):

\[
\begin{pmatrix}
1 + \frac{s_0}{1 + s_0 + t_0}, & \ldots, & 0, & \ldots, & 0, & \ldots, & \frac{s_0}{1 + s_0 + t_0}, & \ldots, & 0, & \ldots, & 0, \\
0, & \ldots, & 0, & \ldots, & 0, & \ldots, & \frac{\hat{s}_1}{1 + \hat{s}_1 + \hat{t}_1}, & \frac{\hat{s}_1}{1 + \hat{s}_1 + \hat{t}_1}, & 0, & \ldots, & 0, \\
\end{pmatrix}
\]
\[
\begin{align*}
(0, 0, 1, 0, \ldots, 0, \ldots, 0, \ldots, 0, \frac{\tilde{t}_0}{1 + s_2 + t_2}, \frac{\tilde{t}_0}{1 + s_2 + t_2}, 0, \ldots, 0) \\
&\vdots \\
(0, 0, 1, 0, \ldots, 0, \ldots, \frac{\tilde{t}_0}{1 + s_0 + t_0}, \frac{\tilde{t}_0}{1 + s_0 + t_0}, 0, \ldots, 0) \quad (3.5) \\
(0, 0, 0, 1, 0, \ldots, 0, \ldots, \frac{\tilde{t}_1}{1 + s_1 + t_1}, \frac{\tilde{t}_1}{1 + s_1 + t_1}, 0, \ldots, 0) \\
&\vdots \\
(0, 0, 0, 0, 1, 0, \ldots, 0, \ldots, 0, \ldots, \frac{\tilde{t}_2}{1 + s_2 + t_2}, \frac{\tilde{t}_2}{1 + s_2 + t_2}, 0, \ldots, 0) \quad (3.6) \\
(0, 0, 0, 0, 0, 1, 0, \ldots, 0, \ldots, 0, \ldots, \frac{\tilde{t}_3}{1 + s_3 + t_3}, \frac{\tilde{t}_3}{1 + s_3 + t_3}, 0, \ldots, 0) \\
&\vdots \\
(0, 0, 0, 0, 0, 0, 1, 0, \ldots, 0, \ldots, 0, 0, \ldots, \frac{\tilde{w}_0}{1 + w_0}, \frac{1}{1 + w_0}, 0, \ldots, 0) \quad (3.7) \\
(0, 0, 0, 0, 0, 0, 0, 1, 0, \ldots, 0, \ldots, 0, 0, 0, \ldots, \frac{1}{1 + w_1}, \frac{1}{1 + w_1}, 0, \ldots, 0) \\
&\vdots \\
(0, 0, 0, 0, 0, 0, 0, 0, 0, \ldots, \frac{1}{1 + w_0}, \frac{1}{1 + w_0}, 0, \ldots, 0) \quad (3.8) \\
&\vdots \\
(0, 0, 0, 0, 0, 0, 0, 0, 0, \ldots, \frac{1}{1 + w_1}, \frac{1}{1 + w_1}, 0, \ldots, 0) \quad (3.9) \\
&\vdots \\
(0, 0, 0, 0, 0, 0, 0, 0, 0, 0, \ldots, \frac{1}{1 + w_2}, \frac{1}{1 + w_2}, 0, \ldots, 0) \\
\end{align*}
\]

First lines of each of the four first sub-sets above, which correspond to the variables \( s_0 \) (5.4), \( t_0 \) (5.5), \( s_0 \) (5.6) and \( t_0 \) (5.7), can be replaced by simpler versions (without changing the determinant) by adding a multiple of the row associated with \( \tilde{w}_0 \) (3.8) or \( w_0 \) (5.9). They become as follows:

\[
\begin{align*}
(1, 0, 0, \ldots, 0, \ldots, 0, \ldots, 0) & (3.6) \\
(0, 1, 0, \ldots, 0, \ldots, 0, \ldots, 0) & (3.7) \\
(0, 0, 1, 0, \ldots, 0, \ldots, 0, \ldots, 0) & (3.9)
\end{align*}
\]
After the above replacement, the \((3f+3r-2) \times (3f+3r-2)\) matrix \(\mathcal{J} = \delta_{jk} - \frac{\eta_{ij}}{\varphi_k \varphi_j}\) acquires the following block-matrix form:

\[
\mathcal{J} = \begin{pmatrix}
1_{(2f+2r) \times (2f+2r)} & B \\
C & D
\end{pmatrix}, \quad \text{whence} \quad J = \det \mathcal{J} = \det(D - CB),
\]

where the matrices \(B, C, D\) are

\[
C = \begin{pmatrix}
\frac{\bar{w}_0}{1+\bar{w}_0}, & \cdots, & 0, & \cdots, & 0, & \cdots, & -\frac{\bar{w}_0}{1+\bar{w}_0}, & \cdots \\
0, & \cdots, & 0, & \cdots, & 0, & \cdots, & 0, & \cdots \\
\vdots & & \ddots & & \vdots & & \ddots & \\
0, & \cdots, & 0, & \cdots, & 0, & \cdots, & 0, & \cdots \\
-\frac{\bar{w}_0}{1+\bar{w}_0}, & \cdots, & 0, & \cdots, & 0, & \cdots, & -\frac{\bar{w}_0}{1+\bar{w}_0}, & \cdots \\
0, & \cdots, & 0, & \cdots, & 0, & \cdots, & 0, & \cdots \\
\vdots & & \ddots & & \vdots & & \ddots & \\
0, & \cdots, & 0, & \cdots, & 0, & \cdots, & 0, & \cdots
\end{pmatrix},
\]
Subtraction of the product \( D \) respectively. The rest of the rows in the matrix \( D \) this process, we shall end up with the expression

\[
\begin{pmatrix}
\frac{1}{1+w_0}, & 0, & 0, \ldots, 0, & 0, \ldots, 0 \\
\frac{1}{1+w_1}, & 0, & 0, \ldots, 0, & 0, \ldots, 0 \\
\vdots & \vdots & \vdots & \vdots & \vdots \\
0, \ldots, -\frac{w_{r-2}}{1+w_{r-2}}, & 1, & 0, \ldots, 0, & 0, \ldots, 0 \\
0, \ldots, 0, & 0, & 0, & -\frac{w_0}{1+w_0}, & 1, \ldots, 0 \\
\vdots & \vdots & \vdots & \vdots & \vdots \\
0, \ldots, 0, & 0, & 0, & 0, \ldots, -\frac{w_{r-2}}{1+w_{r-2}}, & 1, \ldots, 0 \\
\end{pmatrix}.
\quad (3.10)
\]

Therefore, the calculation of the Jacobian has been reduced to finding an explicit formula for the determinant of a smaller, \((\bar{r} + r - 2) \times (\bar{r} + r - 2)\), matrix \( D - CB \). Subtraction of the product \( CB \) does only change two rows of the triangular matrix \( D \), those associated with the variables \( \bar{w}_0 \) and \( w_0 \). They now become

\[
\begin{align*}
\bar{w}_0 &= \frac{1}{1+w_0} \left( 1 + \frac{s_0(1+w_0)}{(1+s_0+t_0)w_0} \right) \\
&\quad + \frac{s_1}{1+s_1+t_1}, \\
&\quad + \frac{s_2}{1+s_2+t_2}, \ldots, \frac{s_{r-2}}{1+s_{r-2}+t_{r-2}} + \frac{s_{r-1}}{1+s_{r-1}+t_{r-1}}, \\
&\quad + \frac{t_0(1+w_0)}{(1+s_0+t_0)w_0} + \frac{t_1}{1+s_1+t_1}, \\
&\quad + \frac{t_2}{1+s_2+t_2}, \ldots, \frac{t_{r-2}}{1+s_{r-2}+t_{r-2}} + \frac{t_{r-1}}{1+s_{r-1}+t_{r-1}},
\end{align*}
\]

respectively. The rest of the rows in the matrix \( D - CB \) are the same as in the matrix \( D \) above, that is they all have only two non-zero elements: one is on the diagonal and the other one is to the left of it. By adding multiples of the columns, without changing the determinant, we can arrange that all those non-diagonal entries in the selected \((\bar{r} - 2) + (r - 2)\) rows vanish, allowing us to compute the determinant by reducing it down to a \( 2 \times 2 \) determinant.

Indeed, multiply the last column of the matrix \( D - CB \) (look at formula \ref{3.10}) by \( w_{r-2} \) and add the result to the penultimate column. The last row has now got only one (diagonal) element. Hence, we keep the factor \( \frac{1}{1+w_{r-2}} \) and reduce the determinant to its minor, removing the last column and the last row. By repeating this process, we shall end up with the expression

\[
J = \frac{\bar{w}_0 w_0 (S\bar{S} - T\bar{T})}{(1+w_0) \cdots (1+w_{r-2})(1+w_0) \cdots (1+w_{r-2})},
\]
where
\[
S = \frac{1}{w_0} + \frac{s_0(1+w_0)}{(1+s_0+t_0)w_0} + \frac{s_1(1+w_1)}{1+s_1+t_1} + \sum_{j=2}^{r-2} \frac{s_j(1+w_j) \prod_{k=1}^{j-1} w_k}{1+s_j+t_j} + \frac{s_{r-1} \prod_{k=1}^{r-2} w_k}{1+s_{r-1}+t_{r-1}},
\]
\[
\tilde{S} = \frac{1}{w_0} + \frac{s_0(1+w_0)}{(1+s_0+t_0)w_0} + \frac{s_1(1+w_1)}{1+s_1+t_1} + \sum_{j=2}^{r-2} \frac{s_j(1+w_j) \prod_{k=1}^{j-1} w_k}{1+s_j+t_j} + \frac{s_{r-1} \prod_{k=1}^{r-2} w_k}{1+s_{r-1}+t_{r-1}},
\]
\[
T = \frac{t_0(1+w_0)}{(1+s_0+t_0)w_0} + \frac{t_1(1+w_1)}{1+s_1+t_1} + \sum_{j=2}^{r-2} \frac{t_j(1+w_j) \prod_{k=1}^{j-1} w_k}{1+s_j+t_j} + \frac{t_{r-1} \prod_{k=1}^{r-2} w_k}{1+s_{r-1}+t_{r-1}},
\]
\[
\tilde{T} = \frac{\tilde{t}_0(1+\tilde{w}_0)}{(1+\tilde{s}_0+\tilde{t}_0)\tilde{w}_0} + \frac{\tilde{t}_1(1+\tilde{w}_1)}{1+\tilde{s}_1+\tilde{t}_1} + \sum_{j=2}^{\tilde{r}-2} \frac{\tilde{t}_j(1+\tilde{w}_j) \prod_{k=1}^{j-1} \tilde{w}_k}{1+\tilde{s}_j+\tilde{t}_j} + \frac{\tilde{t}_{r-1} \prod_{k=1}^{\tilde{r}-2} \tilde{w}_k}{1+\tilde{s}_{r-1}+\tilde{t}_{r-1}}.
\]

This is apparently equivalent to the statement of the Theorem.

4. Hypergeometric series

Now we have all the ingredients for the right-hand side of the Lagrange formula (1.8) and can calculate the expansion coefficients. In this section we return to the initial ordering \( \xi = (x, y, z, \bar{x}, \bar{y}, \bar{z}) \), \( \eta = (s, t, w, \bar{s}, \bar{t}, \bar{w}) \) and \( \varphi = (f, g, h, \bar{f}, \bar{g}, \bar{h}) \).

Let us choose the function \( \chi(\eta) \) in (1.8) to be a Laurent monomial \( \chi(\eta) = \eta^q' = z^m t^n w^p \bar{z}^{\tilde{m}} \bar{t}^{\tilde{n}} \bar{w}^{\tilde{p}}, q' \in \mathbb{Z}^D = \mathbb{Z}^{3d-5} \).

Define the step function \( \sigma_{ij} \) as
\[
\sigma_{ij} \equiv \begin{cases} 
0, & i \leq j \\
1, & i > j 
\end{cases} \quad (4.1)
\]

We shall use the binomial, trinomial and quadrinomial coefficients defined for integer \( m, n, p \) as
\[
\binom{a}{m} \equiv \frac{[x^m](1+x)^a}{m!}, \quad \binom{a}{m, n} \equiv \frac{[x^m y^n](1+x+y)^a}{m!n!}, \quad (4.2a)
\]
\[
\binom{a}{m, n, p} \equiv \frac{[x^m y^n z^p](1+x+y+z)^a}{m!n!p!}, \quad (4.2b)
\]

Note that the multinomial coefficients are evaluated as
\[
\binom{a}{m} = \frac{a(a-1) \ldots (a-m+1)}{m!}, \quad \binom{a}{m, n} = \frac{a(a-1) \ldots (a-m-n+1)}{m!n!}, \quad (4.3a)
\]
\[
\binom{a}{m, n, p} = \frac{a(a-1) \ldots (a-m-n-p+1)}{m!n!p!} \quad (4.3b)
\]

for \( m \geq 0, n \geq 0 \), and \( p \geq 0 \) and vanish if \( m < 0 \), or \( n < 0 \), or \( p < 0 \).

Let \( |m| = m_0 + \ldots + m_{r-1} \) etc., and set
\[
p_{-1} \equiv |m| + |\tilde{n}|, \quad \tilde{p}_{-1} \equiv |\tilde{m}| + |n| \quad (4.4)
\]

We shall also assume that \( p_{r-1} = \tilde{p}_{r-1} = 0 \). In the rest of the paper we shall frequently present formulae for the untilded quantities only, assuming, unless otherwise stated, that the tilded versions are obtained by tilde-symmetry.

Theorem 4.1. The expansion of the monomial \( \eta^q' \) in \( \xi^q \) is given by
\[
\eta^q' = \sum_{q' \geq q} \mathcal{H}_{q'q} \xi^{q'} \quad (4.5)
\]
where

\[ \mathcal{H}_{q'q} = \sum_{i=-1}^{r-1} \sum_{j=-1}^{r-1} \mathcal{F}_{q'q} \mathcal{F}_{q'q}' - \sum_{i=0}^{r-1} \sum_{j=0}^{r-1} \mathcal{G}_{q'q} \mathcal{G}_{q'q}' \] (4.6)

\[ \mathcal{F}_{q'q} = \prod_{j=0}^{r-1} \left( -p_{j-1} - p_j - \delta_{ij} \right) \left( \frac{p_{j-1} + p_j + \delta_{ij} - 1}{p_j - p_j' - \sigma_{ij}} \right), \] (4.7a)

\[ \mathcal{G}_{q'q} = \prod_{j=0}^{r-1} \left( -p_{j-1} - p_j - \delta_{ij} \right) \left( \frac{p_{j-1} + p_j + \delta_{ij} - 1}{p_j - p_j' - \sigma_{ij}} \right), \] (4.7b)

(and respective tilded versions).

**Proof.** Using the definitions [2.21] and [2.22] we get for \( q = (m, n, \tilde{m}, \tilde{n}, \tilde{p}) \):

\[ \varphi^q(\eta) = \frac{f^m g^n h^p f^m \tilde{g}^n h^{\tilde{p}}}{f_0^m \cdots f_{r-1}^m g_0^n \cdots g_{r-1}^n h_0^{r-2} \cdots h_{r-1}^{r-2} \cdots} \]

\[ = \left( \frac{1 + w_0}{1 + s_0 + t_0} \right)^{m+n+p_0} \left( \frac{1 + \tilde{w}_0}{1 + \tilde{s}_0 + \tilde{t}_0} \right)^{\tilde{m}+\tilde{n}+\tilde{p}_0} \]

\[ \times \prod_{j=1}^{r-1} \left( \frac{1 + w_j}{1 + s_j + t_j} \right)^{p_{j-1}+p_j} \prod_{j=1}^{\tilde{r}-1} \left( \frac{1 + \tilde{w}_j}{1 + \tilde{s}_j + \tilde{t}_j} \right)^{\tilde{p}_{j-1}+\tilde{p}_j}, \]

where we assume \( w_{r-1} = \tilde{w}_{\tilde{r}-1} = p_{r-1} = \tilde{p}_{\tilde{r}-1} = 0 \).

Using the shorthand notation [4.3] we can write down \( \varphi^q(\eta) \) in the compact form

\[ \varphi^q(\eta) = \prod_{j=0}^{r-1} \left( \frac{1 + w_j}{1 + s_j + t_j} \right)^{p_{j-1}+p_j} \prod_{j=0}^{\tilde{r}-1} \left( \frac{1 + \tilde{w}_j}{1 + \tilde{s}_j + \tilde{t}_j} \right)^{\tilde{p}_{j-1}+\tilde{p}_j}. \]

Using the step function \( \sigma_{ij} \) [4.1] we rewrite the expressions [4.2] for the ingredients of the Jacobian \( J \) in the following equivalent form:

\[ T = \prod_{i=0}^{r-1} \prod_{j=0}^{r-1} \frac{w_{ij} t_j^{\delta_{ij}}}{(1 + w_j)^{1 - \delta_{ij}} (1 + s_j + t_j)^{\delta_{ij}}}, \]

\[ S = \prod_{i=1}^{r-1} \prod_{j=0}^{r-1} \frac{w_{ij} s_j^{\delta_{ij}}}{(1 + w_j)^{1 - \delta_{ij}} (1 + s_j + t_j)^{\delta_{ij}}}. \]

Substituting the above expressions for \( \chi, \varphi^q \) and \( J \) into the right-hand side of the Lagrange formula [1.3] we get

\[ \chi \varphi^q J = \sum_{i=-1}^{r-1} \sum_{j=-1}^{r-1} \mathcal{F}_i \mathcal{F}_j - \sum_{i=0}^{r-1} \sum_{j=0}^{r-1} \mathcal{G}_i \mathcal{G}_j, \] (4.8)
where

$$F^i = s^{m^i} t^{n^i} w^{p^i} \times \prod_{j=0}^{r-1} \left( \frac{1 + w_j}{1 + s_j + t_j} \right)^{p_{j-1} + p_j} \left( \frac{s_j}{1 + s_j + t_j} \right)^{\delta_{ij}} \frac{w_i^{\sigma_{ij}}}{(1 + w_j)^{1 - \delta_{ij}}}$$

$$= \prod_{j=0}^{r-1} \frac{s_j^{m^i_j + \delta_{ij}} t_j^{n^i_j} w_j^{p^i_j + \sigma_{ij}}}{\prod_{j=0}^{r-1} (1 + s_j + t_j)^{-p_{j-1} - p_j - \delta_{ij}} (1 + w_j)^{-p_{j-1} - p_j - \delta_{ij} - 1}}, \quad (4.9)$$

and \(\tilde{F}^i, \tilde{G}^i\) are their tilded versions. In the above formulae it is assumed that \(w_r^{\sigma_r - 1 - r} = 0 \equiv 1\).

It remains to take the coefficient of the resulting expression for \(\chi \varphi q^J\) at the monomial \(\eta^q\). Using the binomial and trinomial expansions \((4.12)\) and \((4.13)\), the expansion coefficients \(F_{q^i}^j\) and \(G_{q^i}^j\) may be written as a finite sum (the expressions (4.7a) and (4.7b) for the expansion coefficients \(F_{q^i}^j\) and \(G_{q^i}^j\) follow immediately. \)

As formula \((4.16)\) shows, the monomial \(\eta^q\) expands into \(\xi^q\) as a finite sum (the number of the terms equals to the number of the terms in the Jacobian \(J\) of Laurent series. These Laurent series have the uniform structure and can be in fact expressed in terms of a single standard series.

Let us introduce the function \(\Phi(\xi; \mu, \tilde{\mu})\) of \(D = 3d - 5\) complex parameters \(\xi\) and depending on two integer vectors \(\mu = (\mu_0, \ldots, \mu_{r-1}) \in \mathbb{Z}^r\) and \(\tilde{\mu} = (\tilde{\mu}_0, \ldots, \tilde{\mu}_{\tilde{r}-1}) \in \mathbb{Z}^{\tilde{r}}\), altogether \(r + \tilde{r} = d - 1\) integer parameters:

$$\Phi(\xi; \mu, \tilde{\mu}) = \sum_{q \geq 0} (-1)^{|m| + |n| + |\tilde{m}| + |\tilde{n}|} \xi^q$$

$$\times \prod_{j=0}^{r-1} \left( \frac{\mu_j + m_j + n_j + p_{j-1} + p_j - 1}{m_j, n_j, p_j} \right) \prod_{j=0}^{\tilde{r}-1} \left( \frac{\tilde{\mu}_j + \tilde{m}_j + \tilde{n}_j + \tilde{p}_{j-1} + \tilde{p}_j - 1}{\tilde{m}_j, \tilde{n}_j, \tilde{p}_j} \right).$$

(4.11)

Note that, despite the appearance, \(\Phi(\xi; \mu, \tilde{\mu})\) does not factorise into tilde and no-tilde factors because of \(p_{-1} = |m| + |n|\).

For \(\mu_j, \tilde{\mu}_j \geq 1\) the coefficients of the series \((4.11)\) can be written down in terms of the Pochhammer symbols \((a)_m = a(a + 1) \ldots (a + m - 1)\):

$$\Phi(\xi; \mu, \tilde{\mu}) = \sum_{q \geq 0} \frac{(-x)^m (-y)^n z^p (-\bar{x})^\tilde{m} (-\bar{y})^n \bar{z}^\tilde{p}}{m!n!p!|m|!|n|!|p|!}$$

$$\times \prod_{j=0}^{r-1} \left( \frac{\mu_j + m_j + n_j + p_{j-1} + p_j - 1}{(\mu_j)_{p_j}} \right) \prod_{j=0}^{\tilde{r}-1} \left( \frac{\tilde{\mu}_j + \tilde{m}_j + \tilde{n}_j + \tilde{p}_{j-1} + \tilde{p}_j - 1}{(\tilde{\mu}_j)_{\tilde{p}_j}} \right).$$

(4.12)

where \(m! \equiv m_0! \ldots m_{r-1}!\) etc.

Formula \((4.12)\) characterizes \(\Phi\) as a generalised Appel-Horn-type series: the ratios of adjacent coefficients of the power series are rational functions of the indices \(|\mu|, |\tilde{\mu}|\). Note that \(\Phi(0) = 1\), and the series converges in a neighbourhood of 0. The function \(\Phi\) can also be viewed as an \(A\)-hypergeometric function in the sense of
Theorem 4.2. The expansion of the monomial \(\eta^q\) in \(\xi^q\) can be expressed in terms of the function \(\Phi(\xi; \mu, \bar{\mu})\) as follows:

\[
\eta^q = \xi^q \left( \sum_{i=0}^{r-1} \sum_{q=0}^{r-1} \xi^{q^i+\tilde{q}^i} \Phi(\xi; \pi + \nu^i, \bar{\pi} + \bar{\nu}^i) \right)
- \sum_{i=0}^{r-1} \sum_{q=0}^{r-1} \xi^{q_i+\tilde{q}_i} \Phi(\xi; \pi + \nu^i, \bar{\pi} + \bar{\nu}^i) \right),
\]

where

\[
\xi^l = \prod_{j=0}^{r-1} x_j^{\delta_{ij}} \xi_j^{\sigma_{ij}} = \begin{cases} 1, & i = -1 \\ x_0, & i = 0 \\ x_i z_0 \ldots z_{i-1}, & i > 0 \end{cases} \quad (4.14)
\]

\[
\xi^b = \prod_{j=0}^{r-1} y_j^{\delta_{ij}} \xi_j^{\sigma_{ij}} = \begin{cases} 1, & i = -1 \\ y_0, & i = 0 \\ y_i z_0 \ldots z_{i-1}, & i > 0 \end{cases} \quad (4.15)
\]

and respectively for the tilded versions. In (4.10) it is assumed that \(p_{i-1} = |m'| + |\bar{m}'|\).

Proof. Note that the sum over \(q \geq q'\) in (4.10) can be safely replaced with the sum over \(q \in \mathbb{Z}^D\), since the multinomial coefficients in (4.7) vanish for the negative values of the lower indices and thus automatically select the correct limits of summation. Then from (4.5) and (4.6) we can express \(\eta^q\) as the finite sum, each term corresponding to a term in the expansion (4.11) of the Jacobian \(J\):

\[
\eta^q = \sum_{i=-1}^{r-1} \sum_{q=0}^{r-1} F^q_i(\xi) - \sum_{i=0}^{r-1} \sum_{q=0}^{r-1} G^q_i(\xi)
\]

\[
F^q_i(\xi) = \sum_{q \in \mathbb{Z}^D} \overline{F}^q_{i,q} \overline{F}^q_i \xi^q, \quad G^q_i(\xi) = \sum_{q \in \mathbb{Z}^D} \overline{G}^q_{i,q} \overline{F}^q_i \xi^q.
\]

Consider the single term \(F^q_i(\xi)\) of the above sum. From (4.17) it follows that the summation can be restricted to the values of the indices for which the multinomial coefficients do not vanish: \(m_j \geq m_j' + \delta_{ij}, n_j \geq n_j', p_j \geq p_j' + \sigma_{ij}, \) resp. \(\overline{m}', \overline{n}', \) \(\overline{p}', \) \(\overline{\bar{m}}', \) \(\overline{\bar{n}}', \) \(\overline{\bar{p}}'.\) Let us shift the summation indices: \(m_j := m_j + m_j' + \delta_{ij}, n_j := n_j + n_j', p_j := p_j + p_j' + \sigma_{ij},\) etc. so that the summation runs over \(q \geq 0:\)

\[
F^q_i(\xi) = \xi^{q^i} \eta^{q_{i-1}^i} \sum_{q \geq 0} \overline{F}^q_{i,q} \overline{F}^q_i \xi^q,
\]

\[
\overline{F}^q_{i,q} = \prod_{j=0}^{r-1} \left( -p_{j-1} - p_j - p_j' - p_j' - 2\sigma_{i,j-1} \right)_{m_j, n_j} \times \left( p_{j-1} + p_j + p_j' + p_j' + 2\sigma_{i,j-1} \right)_{p_j}.
\]
Applying the identity
\[
\binom{-a}{m, n} \binom{a-1}{p} = (-1)^{m+n} \binom{a+m+n-1}{m, n, p}
\] (4.21)
to (4.20) we get another expression for \( F_{q, q}' \):
\[
F_{q, q}' = (-1)^{|m|+|n|} \prod_{j=0}^{r-1} \binom{m_j + n_j + p_j - 1 + p_j' + p_j' - 1 + 2\sigma_{i,j} - 1}{m_j, n_j, p_j} \] (4.22)

Repeating the same steps for \( \tilde{F}_{q, q}' \) and \( G_{q, q}'(\xi) \) and collecting together the obtained expressions we identify the result as (4.13).

The expansions we are ultimately interested in are those of the monomials in the components \( v_j \) and \( \tilde{v}_j \) of the eigenvector \( V \) (2.2), or of their rescaled versions: \( u_j \) and \( \tilde{u}_j \) (2.8).

**Proposition 4.1.** Let \( u^k = (u_1^k, \ldots, u_r^k), k \in \mathbb{Z}^r \), and resp. for \( \tilde{u}^\tilde{k} \). The expansion of the monomial \( u^k \tilde{u}^\tilde{k} \) in \( \xi^q \) is given by the formula
\[
\sum_{i=0}^{r-1} \sum_{j=0}^{\tilde{r}-1} \xi^i \tilde{\Phi}(\xi; \pi + \nu^i, \tilde{\pi} + \tilde{\nu}^i)
- \sum_{i=0}^{r-1} \sum_{j=0}^{\tilde{r}-1} \xi^i \tilde{\Phi}(\xi; \pi + \nu^i, \tilde{\pi} + \tilde{\nu}^i),
\] (4.23)
where \( \Phi, \xi^i, \xi^\tilde{i}, \nu^i \) are same as in Theorem 4.2 and
\[
\pi_j = k_{j+1} + \ldots + k_r, \quad j = 0, \ldots, r - 1,
\] (4.24)

**Proof.** Using formulae (2.14) we can express \( u_j \) as monomials in \( x_0, s_0, z \) and \( w \). Thus
\[
u^j = x^{-m} z^{-p} s^{m} w^{p},
\] (4.25)
where
\[
m_0' = k_1 + k_3 + \ldots, \quad m' = \ldots = m_{r-1}' = 0,
\] (4.26)
\[
p_{2j}' = k_{2j+2} + k_{2j+4} + \ldots, \quad p_{2j}' = k_{2j+3} + k_{2j+5} + \ldots
\] (4.27)
(similarly for the tilded variables). Substituting these values of \( p_j' \) into the formula \( \pi_j = p_j' + p_{j+1}' (4.10) \) we get (4.24). Note that \( p_{r-1}' \equiv |m'| = m_0' \). It remains to apply Theorem 4.2 to the monomial (4.25).

To obtain the expansion of the monomial \( v^k \tilde{v}^\tilde{k} \) in \( bcbc \) one needs to express the components \( v_j \) of the eigenvector \( V \) in terms of \( uu \) using (2.8), then apply Proposition 4.1 and finally, using (2.14) and (2.11), to express \( \xi \) in terms of \( bcbc \). To get the expansions involving the eigenvalue \( \lambda \) one has to use (2.9). In particular, the resulting expansions for \( u_i \) or \( \tilde{u}_i \) give an explicit solution for iterations of equations (2.12).

In this paper we study only formal power series expansions. The analytic version of the Lagrange theorem \( \S \) guarantees that our power series converge in an open neighbourhood around \( \xi = 0 \). The precise description of the convergence domain of a multivariate power series is, however, a difficult task, see for example \( \S 10 \), and we leave it to further study.
5. Corner eigenvalue

As mentioned in the end of section 2, the cases of the corner eigenvalue \( r = 0 \) \((\bar{r} = 0)\) have some peculiarities. Let us consider the case \( \bar{r} = 0 \) in more detail (the case \( r = 0 \) can be treated by the tilde-symmetry). We have

\[
\begin{pmatrix}
0 & b_0 & b_1 \\
c_0 & a_1 & b_1 \\
\vdots & \vdots & \vdots \\
c_{i-1} & a_i & b_i \\
c_{r-1} & a_r & \end{pmatrix} \begin{pmatrix}
v_1 \\
v_2 \\
\vdots \\
v_i \\
v_r 
\end{pmatrix} = \lambda \begin{pmatrix}
v_1 \\
v_2 \\
\vdots \\
v_i \\
v_r 
\end{pmatrix}. \quad (5.1)
\]

The tilded variables are absent. Besides, the sequences \( y_i \), resp. \( t_i \) are also absent. The total set of variables \( \xi = (x, z) \) has the cardinality \( (r)+(r-1) = 2r-1 = 2d-3 \), and all the expansion variables are independent, in contrast with the generic case. Respectively, \( \eta = (s, w) \) and \( \varphi = (f, h) \), where

\[
f_0 = \ldots = f_{r-1} = \frac{1 + w_0}{1 + s_0}, \quad h_i = \frac{(1 + w_i)(1 + w_{i+1})}{(1 + s_i)(1 + s_{i+1})}, \quad i = 0, \ldots, r - 2. \quad (5.2)
\]

The expression (4.1) for the jacobian \( J \) simplifies to

\[
J = S = \left( 1 + \sum_{j=0}^{r-1} s_j (1 + w_j) \prod_{k=j}^{j-1} w_k \right) \prod_{k=0}^{r-2} (1 + w_k)^{-1} \quad (5.3)
\]

and contains thus only \( r + 1 = d \) terms. The analog of formula (4.8) is

\[
\chi \varphi^q J = \sum_{i=-1}^{r-1} F^i, \quad (5.4)
\]

where

\[
F^i = \prod_{j=0}^{r-1} s_j^{m_j + \delta_{ij}} w_j^{p_j + \sigma_{ij}} (1 + s_j)^{-1} \prod_{k=j}^{j-1} (1 + w_j)^{p_k + \delta_{ik} - 1}. \quad (5.5)
\]

The expression (4.6) for the expansion coefficient becomes

\[
H_{q'q} \equiv [\xi^q \eta^{q'} \equiv [x^m z^p]s^{m'} w^{p'} = \sum_{i=-1}^{r-1} F_{q'q}^i, \quad (5.6)
\]

where

\[
F_{q'q}^i = [s^{m} w^{p}] F^i = \prod_{j=0}^{r-1} \left( -p_{j-1} - p_j - \delta_{ij} \right) \prod_{j=0}^{r-2} \left( p_{j-1} + p_j + \delta_{ij} - 1 \right), \quad (5.7)
\]

with \( p_{r-1} = |m| \).

The expansion of the monomial \( u^k \) in \( x^m z^p \) is given by the analog of formula (4.12):

\[
u^k = \sum_{i=-1}^{r-1} \xi^i \Phi(x, z; \pi + \nu^i) \quad (5.8)
\]

where \( \xi^i, \nu^i \) are same as in Theorem (4.2), \( \pi \) is given by (4.24), and \( \Phi(x, z; \mu) \) is defined as the series

\[
\Phi(x, z; \mu) = \sum_{m,p \geq 0} (-x)^m z^p \prod_{j=0}^{r-1} (\mu_j + m_j + p_{j-1} + p_j - 1), \quad (5.9)
\]
or, for \( \mu_j \geq 1 \),
\[
\Phi(x, z; \mu) = \sum_{m,p \geq 0} \frac{(-x)^m z^p}{m!p!} \prod_{j=0}^{r-1} \frac{(\mu_j + p_j + p_{j-1})}{(\mu_j)_{p_j-1}}.
\] (5.10)

6. Examples

In this section we illustrate our general results with a few low-dimensional examples.

6.1 The simplest case is \( d = 2 (r = 1, \tilde{r} = 0) \). The spectral problem is
\[
\begin{pmatrix}
0 & b_0 \\
c_0 & a_1
\end{pmatrix}
\begin{pmatrix}
1 \\
v_1
\end{pmatrix}
= \lambda
\begin{pmatrix}
1 \\
v_1
\end{pmatrix},
\] (6.1)

After eliminating the eigenvalue \( \lambda = b_0 v_1 \) and introducing the rescaled variables \( u \equiv u_1 = -v_1 a_1 / c_0 \), \( x \equiv x_0 = b_0 c_0 / a_1^2 \) we get for the single unknown variable \( u \) the single quadratic equation
\[
1 - u - xu^2 = 0.
\] (6.2)

The branch we are studying is selected by the condition \( u \to 1 \) for \( x \to 0 \), and the corresponding solution is
\[
u(x) = \frac{\sqrt{1 + 4x} - 1}{2x}.
\] (6.3)

In terms of the variable \( s \equiv xu \) we get the equivalent Lagrange-type equation
\[
x = s(1 + s) = \frac{s}{f}, \quad f = \frac{1}{1 + s}.
\] (6.4)

The corresponding Jacobian \( J \) has two terms:
\[
J = 1 - s \frac{df}{ds} = 1 + \frac{s}{1 + s}.
\] (6.5)

The expansion of the monomial \( u^k, k \in \mathbb{Z} \) given by Proposition 4.1 is
\[
u^k = \Phi(x; k) + x \Phi(x; k + 2),
\] (6.6)

where
\[
\Phi(x; \mu) = \sum_{m \geq 0} (-x)^m \frac{(\mu + 2m - 1)}{m!} = \frac{1}{\sqrt{1 + 4x}} \left( \frac{\sqrt{1 + 4x} - 1}{2x} \right)^{\mu - 1},
\] (6.7)

For \( \mu \geq 1 \) the function \( \Phi(x; \mu) \) can be expressed in terms of the Gauss hypergeometric series:
\[
\Phi(x; \mu) = \sum_{m \geq 0} \frac{(-x)^m}{m!} \frac{(\mu)_{2m}}{(\mu)_m} = \text{$_2F_1$} \left( \frac{\mu}{2}, \frac{\mu + 1}{2} ; -4x \right).
\] (6.8)

As a matter of fact, for \( k \geq 0 \) the two hypergeometric series in (6.6) can be summed into a single series
\[
u^k = \sum_{m \geq 0} \frac{(-x)^m}{m!} \frac{(k)_{2m}}{(k + 1)_m} = \text{$_2F_1$} \left( \frac{k}{2}, \frac{k + 1}{2} ; -4x \right) = \left( \frac{\sqrt{1 + 4x} - 1}{2x} \right)^k.
\] (6.9)

6.2 Three-dimensional case, corner eigenvalue: \( d = 3 (r = 2, \tilde{r} = 0) \). The spectral problem
\[
\begin{pmatrix}
0 & b_0 & 0 \\
c_0 & a_1 & b_1 \\
0 & c_1 & a_2
\end{pmatrix}
\begin{pmatrix}
1 \\
v_1 \\
v_2
\end{pmatrix}
= \lambda
\begin{pmatrix}
1 \\
v_1 \\
v_2
\end{pmatrix},
\] (6.10)
after the substitutions \( \lambda = b_0 v_1, v_1 = -u_1 c_0 / a_1, v_2 = u_2 c_0 c_1 / a_1 a_2 \) is reduced to the pair of quadratic equations for the variables \( u_1, u_2 \)

\[
\begin{align*}
1 - u_1 - x_0 u_1^2 + z_0 u_2 & = 0, \\
u_1 - u_2 - x_1 u_1 u_2 & = 0,
\end{align*}
\]

(6.11a) (6.11b)

where the expansion variables are

\[
x_0 = \frac{b_0 c_0}{a_1}, \quad x_1 = \frac{b_0 c_0}{a_1 a_2}, \quad z_0 = \frac{b_1 c_1}{a_1 a_2},
\]

(6.12)

and we choose the branch \( u_1, u_2 \to 1 \) as \( x_0, x_1, z_0 \to 0 \).

The equivalent Lagrange-type equations for the variables \( s_0 = x_0 u_1, s_1 = x_1 u_1, w_0 = z_0 u_2 \) are

\[
x_0 = \frac{s_0}{f_0}, \quad x_1 = \frac{s_1}{f_1}, \quad z_0 = \frac{w_0}{h_0},
\]

(6.13)

where

\[
f_0 = f_1 = \frac{1 + w_0}{1 + s_0}, \quad h_0 = \frac{1 + w_0}{(1 + s_0)(1 + s_1)}.
\]

(6.14)

The three-term Jacobian \( J \) equals

\[
J = \frac{1}{1 + w_0} + \frac{s_0}{1 + s_0} + \frac{s_1 w_0}{(1 + s_1)(1 + w_0)}.
\]

(6.15)

By formula (6.8), the monomial \( u_1^{k_1} u_2^{k_2} \) expands in \( x_0, x_1, z_0 \) as

\[
u_1^{k_1} u_2^{k_2} = \Phi(x_0, x_1, z_0; k_1 + k_2, k_2) + x_0 \Phi(x_0, x_1, z_0; k_1 + k_2 + 2, k_2) + x_1 z_0 \Phi(x_0, x_1, z_0; k_1 + k_2 + 2, k_2 + 2),
\]

(6.16)

where

\[
\Phi(x_0, x_1, z_0; \mu_0, \mu_1) = \sum_{m_0, m_1, p_0 \geq 0} (-x_0)^{m_0} (-x_1)^{m_1} z_0^{p_1} \left( \frac{\mu_0 + 2m_0 + m_1 + p_0 - 1}{m_0, p_0} \right) \left( \frac{\mu_1 + m_1 + p_0 - 1}{m_1} \right),
\]

(6.17)

or, for \( \mu_0, \mu_1 \geq 1 \),

\[
\Phi(x_0, x_1, z_0; \mu_0, \mu_1) = \sum_{m_0, m_1, p_0 \geq 0} (-x_0)^{m_0} (-x_1)^{m_1} z_0^{p_1} \frac{(\mu_0)_{2m_0 + m_1 + p_0} (\mu_1)_{m_1 + p_0}}{(m_0)! m_1! p_0!}.
\]

(6.18)

### 6.3 Three-dimensional case, middle eigenvalue: \( d = 3 (r = \bar{r} = 1) \)

The spectral problem

\[
\begin{pmatrix}
\bar{a}_1 & \bar{b}_0 & 0 \\
\bar{c}_0 & 0 & \bar{b}_0 \\
0 & \bar{c}_0 & \bar{a}_1
\end{pmatrix}
\begin{pmatrix}
\bar{v}_1 \\
v_1 \\
y_1
\end{pmatrix}
= \lambda
\begin{pmatrix}
\bar{v}_1 \\
v_1 \\
y_1
\end{pmatrix}
\]

(6.19)

after the substitutions \( \lambda = \bar{c}_0 \bar{v}_1 + b_0 v_1, v_1 = -u c_0 / a_1, \bar{v}_1 = -\bar{u} \bar{b}_0 / \bar{a}_1 \) is reduced to the pair of quadratic equations for the variables \( u, \bar{u} \)

\[
\begin{align*}
1 - u - x u^2 - y u \bar{u} & = 0, \\
1 - \bar{u} - \bar{x} \bar{u}^2 - \bar{y} u \bar{u} & = 0,
\end{align*}
\]

(6.20a) (6.20b)

where the expansion variables are

\[
x = \frac{b_0 c_0}{a_1^2}, \quad y = \frac{b_0 c_0}{a_1 a_1}, \quad \bar{x} = \frac{\bar{b}_0 \bar{c}_0}{a_1^2}, \quad \bar{y} = \frac{b_0 c_0}{a_1 a_1},
\]

(6.21)

and the chosen branch is \( u, \bar{u} \to 1 \) as \( x, y, \bar{x}, \bar{y} \to 0 \). Note that the variables \( x, y, \bar{x}, \bar{y} \) are bound by the single relation \( \bar{x} \bar{x} = y \bar{y} \).
The equivalent Lagrange-type equations for the variables \( s = xu, \tilde{s} = \tilde{xu}, t = y\tilde{u}, \tilde{t} = \tilde{y}\tilde{u} \) are

\[
x = \frac{s}{f}, \quad y = \frac{t}{g}, \quad \tilde{x} = \frac{\tilde{s}}{\tilde{f}}, \quad \tilde{y} = \frac{\tilde{t}}{\tilde{g}},
\]

where

\[
f = \tilde{g} = \frac{1}{1 + s + t}, \quad \tilde{f} = g = \frac{1}{1 + \tilde{s} + \tilde{t}}.
\]

The five-term Jacobian \( J \) equals

\[
J = 1 + \frac{s}{1 + s + t} + \frac{\tilde{s}}{1 + \tilde{s} + \tilde{t}} + \frac{s\tilde{s}}{(1 + s + t)(1 + \tilde{s} + \tilde{t})} - \frac{\tilde{t}t}{(1 + s + t)(1 + \tilde{s} + \tilde{t})}.
\]

By formula (6.23), the monomial \( u^k\tilde{u}^{\tilde{k}} \) expands in \( \xi = (x, y, \tilde{x}, \tilde{y}) \) as

\[
u^k\tilde{u}^{\tilde{k}} = \Phi(\xi; k, \tilde{k}) + x\Phi(\xi; k + 2, \tilde{k}) + \tilde{x}\Phi(\xi; k, \tilde{k} + 2) + x\tilde{x}\Phi(\xi; k + 2, \tilde{k} + 2) - y\tilde{y}\Phi(\xi; k + 2, \tilde{k} + 2),
\]

where, for \( q = (m, n, \tilde{m}, \tilde{n}) \),

\[
\Phi(\xi; \mu, \tilde{\mu}) = \sum_{q \geq 0} \frac{(-x)^m(-y)^n(-\tilde{x})^{\tilde{m}}(-\tilde{y})^{\tilde{n}}}{m!n!\tilde{m}!\tilde{n}!} \left( \frac{\mu + 2m + n + \tilde{n} - 1}{m}, \frac{\tilde{\mu} + 2\tilde{m} + n + \tilde{n} - 1}{\tilde{m}}, \frac{\mu + 2\tilde{m} + n + \tilde{n} - 1}{\tilde{n}} \right),
\]

or, for \( \mu, \tilde{\mu} \geq 1 \),

\[
\Phi(\xi; \mu, \tilde{\mu}) = \sum_{q \geq 0} \frac{(-x)^m(-y)^n(-\tilde{x})^{\tilde{m}}(-\tilde{y})^{\tilde{n}}}{m!m!\tilde{m}!\tilde{n}!} \frac{(\mu)_{2m+n+\tilde{n}}(\tilde{\mu})_{2\tilde{m}+n+\tilde{n}}}{(\mu)_{m+n}(\tilde{\mu})_{m+n}}.
\]

7. Discussion

In the present paper we have solved a problem of much physical interest. The spectra of finite Jacobi matrices appear in many applications: from orthogonal polynomials and nearest-neighbours interaction models to solvable models of quantum mechanics (Lamé polynomials and Bethe ansatz).

What is left for further study are the questions of convergence domains, differential equations and integral representations for the obtained hypergeometric series, as well as their relation to \( \mathcal{A} \)-hypergeometric functions introduced in [4, 5, 6]. The approach used in our work can be generalised to multiparameter spectral problems [11].

Explicit perturbative solutions have also appeared in a different context in the works by Edwin Langmann [12, 13], notably for the multi-dimensional spectral problems related to the Calogero-Sutherland and elliptic Calogero-Moser systems.

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