Research Article

An Improved Lattice Hydrodynamic Model by considering the Effect of “Backward-Looking” and Anticipation Behavior

Jin Wan, Xin Huang, Wenzhi Qin, Xiuge Gu, and Min Zhao

School of Science, Nantong University, Nantong, Jiangsu 226019, China

Correspondence should be addressed to Min Zhao; zhao.m@ntu.edu.cn

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In order to prevent the occurrence of traffic accidents, drivers always focus on the running conditions of the preceding and rear vehicles to change their driving behavior. By taking into the “backward-looking” effect and the driver’s anticipation effect of flux difference consideration at the same time, a novel two-lane lattice hydrodynamic model is proposed to reveal driving characteristics. The corresponding stability conditions are derived through a linear stability analysis. Then, the nonlinear theory is also applied to derive the mKdV equation describing traffic congestion near the critical point. Linear and nonlinear analyses of the proposed model show that how the “backward-looking” effect and the driver’s anticipation behavior comprehensively affect the traffic flow stability. The results show that the positive constant $c$, the driver’s anticipation time $\tau$, and the sensitivity coefficient $p$ play significant roles in the improvement of traffic flow stability and the alleviation of traffic congestion. Furthermore, the effectiveness of linear stability analysis and nonlinear analysis results is demonstrated by numerical simulations.

1. Introduction

With the rapid development of wireless technologies and intelligent vehicles in the past years, the vehicle-to-vehicle (V2V) communication [1, 2] can not only alleviate the congested traffic flow but also provide pleasurable driving experiences. To reveal the realistic traffic phenomena, there have emerged various traffic models such as car-following models [3–10], cellular automation models [11, 12], macro-traffic models [13–15], lattice hydrodynamic models [16–21], continuum models [22–24], and gas kinetic models [25].

By incorporating the ideas of car-following as well as continuum models, the first lattice hydrodynamic model [26] with a conservation equation (1) and a flow equation (2) was proposed in 1998.

$$\begin{align*}
\partial_t \rho_j + \rho_0 \left( \rho_j v_j - \rho_{j-1} v_{j-1} \right) &= 0, \\
\partial_t \left( \rho_j v_j \right) &= a \left[ \rho_0 V \left( \rho_{j+1} \right) - \rho_j v_j \right],
\end{align*}$$

where $\rho_0 V \left( \rho_{j+1} \right)$ and $\rho_j v_j$ denote the optimal flow and the actual flow, respectively. $\rho_0$ is the average density, and $a$ refers to the driver’s sensitivity. Then, the stability condition and mKdV equation can be, respectively, obtained with the linear analysis and nonlinear analysis methods. Subsequently, many extended works have been developed by taking different factors into accounts, such as driver’s memory [27–31], driver’s anticipation effect [32–34], density difference [35], traffic interruption probability [36–38], and “backward-looking” effect [39–41].

In addition to the information from the preceding vehicle, drivers always focus on the running conditions of the rear vehicles from the rear-view mirror to change their driving behavior. By considering the “backward-looking” effect, Ge and Cheng [39] presented the BL-LV model as follows:

$$\begin{align*}
\rho_j (t + \tau) v_j (t + \tau) &= \rho_0 V \left( \rho_{j+1} (t) \right) \\
&+ H(p - \rho_{j+1} (t)) \cdot H(\rho_{j-1} (t) - \rho_c) \rho_0 V_f \left( \rho_{j-1} (t) \right),
\end{align*}$$

where $\rho_0 V \left( \rho_{j+1} \right)$ and $\rho_j v_j$ denote the optimal flow and the actual flow, respectively. $\rho_0$ is the average density, and $a$ refers to the driver’s sensitivity. Then, the stability condition and mKdV equation can be, respectively, obtained with the linear analysis and nonlinear analysis methods. Subsequently, many extended works have been developed by taking different factors into accounts, such as driver’s memory [27–31], driver’s anticipation effect [32–34], density difference [35], traffic interruption probability [36–38], and “backward-looking” effect [39–41].

In addition to the information from the preceding vehicle, drivers always focus on the running conditions of the rear vehicles from the rear-view mirror to change their driving behavior. By considering the “backward-looking” effect, Ge and Cheng [39] presented the BL-LV model as follows:
where $H(\cdot)$ and $\rho_c$ represent the Heaviside function and the safety density, respectively. $\bar{\rho}$ denotes a parameter close to 1. $V_p(\cdot)$ is the optimal velocity function for "forward-looking." While $V_f(\cdot)$ is the optimal velocity function for "backward-looking," which works if and only if $\rho_c \leq \rho - \bar{\rho} \leq \bar{\rho}$ holds. Functions $V_p(\cdot)$ and $V_f(\cdot)$ have exactly the opposite effects and are defined as

$$V_p(\rho_{j+1}) = \tanh\left(\frac{2}{\rho_0 - \rho_{j+1}} - \frac{1}{\rho_c}\right) + \tanh\left(\frac{1}{\rho_c}\right),$$

$$V_f(\rho_{j-1}) = \gamma\left[-\tanh\left(\frac{2}{\rho_0 - \rho_{j-1}} - \frac{1}{\rho_c}\right) + \tanh\left(\frac{1}{\rho_c}\right)\right],$$

where $\gamma$ is a positive constant that represents the relative role of the "backward-looking." For the sake of safe and comfortable driving, drivers always predict the velocity of the preceding vehicles so that they can adjust their driving velocity at the next moment. This driving behavior is known as the anticipation effect [42] or predictive effect [43, 44]. Wang et al. [43] added the driver’s anticipation effect into the optimal velocity (OV) function in a single-lane lattice model. Later, Kaur and Sharma [44] extended the single-lane model [43] to the two-lane case along with the optimal current difference (OCD) effect. By considering the anticipation effects in the OCD term as well as the OV function, Sharma [45] constructed a new lattice hydrodynamic model with an anticipation coefficient $\alpha$, where the positive $\alpha$ represents the anticipation effect while the negative $\alpha$ represents the delay effect. Then, Li et al. [32] adopted this anticipation effect to the multiple optimal current differences’ anticipation version. Meanwhile, numerical results indicated that only the information of three preceding lattices [46] was enough for the improvement of traffic stability. In view of the anticipation effect of flux difference, Chang and Cheng [47] proposed a new lattice model as follows:

$$\partial_t(\rho_j v_j) = \rho_0 V(\rho_{j+1}) - a(1 + k_1) \rho_j v_j + k_2(\rho_{j+1}(t + \tau)v_{j+1}(t + \tau) - \rho_j v_j),$$

where $p$ is the sensitivity coefficient of the anticipation flux difference. The greater $p$ is, the stronger the driver’s anticipation effect is. $\tau$ means the anticipation time, and $\rho_{j+1}(t + \tau)v_{j+1}(t + \tau) - \rho_j v_j$ represents the anticipation flux difference between the preceding lattice $j + 1$ and the current lattice $j$. The optimal velocity functions and the function $H(\cdot)$ are the same as those in Ge and Cheng’s model [39].

After eliminating the velocity term $v_j$ in equations (1) and (6), the following density equation can be derived:
When $p = 0$, the model reduces to Ge and Cheng’s model [39]. In addition, the lattice model coincides with Nagatani’s model [26] when $p = 0$ and $\tau = 0$.

The influence of the “backward-looking” effect and the anticipation effect has been studied by linear stability analysis. The uniform traffic stream with the fixed density $\rho_0$ and the optimal velocity $V'(\rho_0)$ is defined as the stable state. One can obtain the solution of the steady-state vehicular system as follows:

$$\rho_j(t) = \rho_0,$$

$$\nu(j) = V_F(\rho) + V_B(\rho).$$

Then, $y_j$ is supposed as a small perturbation added into the steady-state solution at lattice $j$, and the perturbed solution is

$$\rho_j(t) = \rho_0 + y_j(t).$$

Equation (9) is substituted into equation (7), and the derived equation is linearized. The equation is obtained as

$$\begin{align*}
\partial_t^2 y_j(t) + a\partial_t y_j(t) &+ (a\rho^2 \frac{\partial^2}{\partial x^2} V'(\rho_0))(y_{j+1}(t) - y_{j-1}(t)) \\
+ a\partial_t y_j &+ \nu(j_0) (y_{j+1}(t) - y_{j-1}(t)) \\
- a\partial_t y_j &+ y_j(t) + (a + a)\partial_t y_j = 0.
\end{align*}$$

Let the small perturbation $y_j(t) = \exp(i k x + z t)$, and inserting it into equation (10), we gain

$$\begin{align*}
\partial_t^2 z &+ a^2 \rho_0^2 (V_F'(\rho_0)(\exp(i k) - 1) + V_B'(\rho_0)(\exp(i k) - \exp(-i k))) \\
+ a\partial_t z &+ (a + a)z = 0.
\end{align*}$$

Assuming $z = z_1(ik) + z_2(ik)^2 + \cdots$ and putting it into equation (11), we reserve the first-order and second-order terms of $ik$ by neglecting higher-order terms. Then, one can get

$$\begin{align*}
z_1 &= -\rho_0^2 (V_F'(\rho_0) + V_B'(\rho_0)), \\
z_2 &= \frac{2a\rho z_1 - 2z_1^2 + 2a\rho z_1^2 - 2(aV_F'(\rho_0) + 3V_B'(\rho_0)\rho_0^2 + 3aV_B'(\rho_0)\rho_0^2)}{2a}.
\end{align*}$$

The uniform steady-state flow remains stable as $z_2$ is positive. On the contrary, the uniform flow tends to be unstable if the value of $z_2$ is negative. When $z_2$ is equal to zero, the formula of the neutral stability condition is acquired as

$$a > \frac{2(V_F'(\rho_0) + V_B'(\rho_0))^2 \rho_0^2}{M}$$

It is obvious that the stable condition (14) is reduced to that of Ge and Cheng’s model [39] when $p = 0$. Equation (14) shows that the parameters $r$, $\rho$, and $\gamma$ play important roles in the performance of the traffic system.

The phase diagram of $(\rho, a)$ is shown in Figure 1, where $\rho$ is density and $a$ denotes sensitivity. Note that $\gamma$ denotes a positive constant considering the “backward-looking” effect [40, 41]. Three solid curves in each pattern of Figure 1 represent the neutral stability curves under $\gamma = 0.1$. The phase diagram is divided into two regions: the stable and unstable regions are above and below the solid neutral stability curve, respectively. In the stable region, the traffic flow will remain stable with a perturbation. In the unstable region, the congested traffic will happen even with a small perturbation. Figure 1 depicts the neutral stability curves of the FADBE model for different values of driver’s anticipation time $\tau$ and the sensitivity coefficient $p$ of the anticipation flux difference when $\gamma = 0.1$. As shown in Figure 1, the corresponding stable region expands gradually with the increasing values of $\tau$ and $p$, which means that the driver’s anticipation effect can contribute to the stabilization of traffic flow obviously.

Figure 2 demonstrates the neutral stability curves of the FADBE model for different values of $\tau$ and $\gamma$ when $\rho = 0.1$. As shown in Figure 2, the corresponding neutral stability curves and critical points decrease gradually with the increasing values of $\tau$ and $\gamma$. When the value of $\gamma$ is fixed in each pattern of Figure 2, the stability of the traffic flow can be improved efficiently with the increasing value of anticipation time $\tau$.

Figure 3 exhibits that the phase diagram of $(\rho, a)$ of the FADBE model for different values of the positive constant $\gamma$ ($\gamma = 0.05, 0.1, 0.15, 0.2$) and the sensitivity coefficient $p$ ($p = 0.1, 0.2, 0.3$) when $\tau = 1.5$. With the increasing value of the positive constant $\gamma$ in each pattern of Figure 3, the amplitudes of neutral stability curves all fall down gradually when the value of the parameter $p$ is fixed. For different values of $\gamma$, it can be seen from each pattern of Figure 3 that the corresponding peak value of solid curves lowers down with the increasing value of $\rho$. These phenomena demonstrate that when the value of $\gamma$ or $p$ increases, the area of the stable region becomes larger. Thus, it can be concluded that the traffic flow stability can be improved by incorporating the positive constant $\gamma$ and the sensitivity coefficient $p$.

### 3. Nonlinear Stability Analysis

To investigate the influence of the “backward-looking” effect and the anticipation effect of flux difference, we carried out the reduction perturbation method to obtain the mKdV equation. The slow variables $X$ and $T$ are defined as follows:
\[ X = \varepsilon (j + bt), \]
\[ T = \varepsilon^3 t, \]
\[ \rho_j = \rho_c + \varepsilon R(X, T), \]

where \( 0 < \varepsilon \ll 1 \) is a small positive scaling parameter. Substituting equation (15) into equation (7), each term of equation (7) is expanded up to the fifth-order of \( \varepsilon \) with the Taylor expansion technique as follows:
\begin{align}
&\varepsilon^3 \left[ b + \rho \left( V_F' + V_B' \right) \right] \partial_X R \\
&+ \varepsilon^3 \left( \frac{b^2}{a} - b^3 \rho - \frac{3}{2} \partial^4 \rho - \frac{1}{2} \partial^2 V_F' \right) \partial_X R \\
&+ \varepsilon^4 \left[ \partial_X R \left( \frac{1}{2} b^3 \rho - \frac{7}{2} \partial^2 \rho - \frac{1}{6} \partial^3 V_F' \right) \partial_X R + \left( \frac{1}{6} \partial V_B' + \frac{1}{6} \partial^2 V_F' \right) \partial_X R \right] \\
&+ \varepsilon^5 \left[ \left( \frac{2b}{a} - 2bp \right) \partial_X \partial_X R + \left( -\frac{1}{4} b^4 \rho - \frac{5}{8} \partial^4 \rho + \frac{1}{24} \partial^2 V_F' \right) \partial_X R + \left( \frac{1}{24} \partial V_B' + \frac{1}{24} \partial^2 V_F' \right) \partial_X R \right] = 0,
\end{align}

\textit{Figure 2:} The neutral stability curves for $p = 0.1$ for different values of $\tau$ and $\gamma$: (a) $\tau = 0.5$; (b) $\tau = 1.5$; (c) $\tau = 2.5$; (d) $\tau = 3.5$. 

\textit{Complexity 5}
where
\[ V'_F = (dV/d\rho)|_{p=p_c} \]
and
\[ V'_B = (dV/d\rho)|_{p=p_c} \]  
\[ V^2_B = (d^3V/d\rho^3)|_{p=p_c} \] . Near the critical point \((\rho_c, a_c)\), the value of \(a_c\) is defined as \(a_c = a(1 + \varepsilon^2)\). The squared and cubic terms of \(\varepsilon\) are eliminated in equation (16) with the consideration of \(b = -\rho_c^2 (V'_F + V'_B)\). One can get the simplified equation as follows:
\[
\varepsilon^4 \left[ \partial_\tau R - g_1 \partial_X^3 R + g_2 \partial_X^5 R \right] 
+ \varepsilon^5 \left[ g_3 \partial_X^2 R + g_4 \partial_X^4 R + g_5 \partial_X^2 R^3 \right] = 0, 
\] (17)
where

\[ \partial_\tau \]
The solution of the mKdV equation can be rewritten as transformations are applied:

\[ g_1 = \frac{1}{6} \rho_c \left( V'_b + V'_f + 3 \rho_p \varepsilon^2 (V'_b + V'_f)^3 \right), \]

\[ g_2 = \rho_c^2 \left( V'_b + V'_f \right), \]

\[ g_3 = \frac{1}{2} \rho_c^2 \left( 3V'_b - V'_f - 2 \rho (V'_b + V'_f) + 2 \rho_p \varepsilon^2 (V'_b + V'_f)^3 \right), \]

\[ g_4 = -\frac{1}{24} \rho_c^2 \left( 15V'_b - V'_f + 4 \rho_p \varepsilon^2 (V'_b + V'_f)^3 \right) - \frac{1}{3a_c^4} \left( a \rho \right)(a \rho - 1)(V'_b + V'_f) \left( 3 \rho_p \varepsilon^2 (V'_b)^3 + V'_f \right) \]

\[ -\frac{1}{3a_c^4} \left( a \rho \right)(a \rho - 1)(V'_b + V'_f) \left( 3 \rho_p \varepsilon^2 (V'_b)^3 + V'_f \right), \]

\[ g_5 = -\frac{1}{12} \rho_c^2 \left( 3V'_b - V'_f \right) - \frac{1}{3a_c^4} \left( a \rho \right)(a \rho - 1)(V'_b + V'_f) \left( V''_b + V''_f \right). \]

To derive the standard mKdV equation, the following transformations are applied:

\[ T = \frac{1}{\rho_c} T', \]

\[ R = \frac{g_1}{g_2}. \]

Thus, equation (17) is converted into the standard mKdV equation as follows:

\[ \partial_T R' = \partial_x^3 R' - \partial_x R^3 + \varepsilon \left[ \frac{g_1}{g_2} \partial_x^2 R' + \frac{g_2}{g_1} \partial_x^2 R' + \frac{g_3}{g_2} \partial_x^2 R^3 \right]. \]

(20)

When the term \( O(\varepsilon) \) is ignored, the kink-antikink solution of the mKdV equation can be rewritten as

\[ R'_0(X, T') = \sqrt{c} \tanh \left( \sqrt{c} \left( X - cT' \right) \right), \]

(21)

where \( c \) represents the propagation velocity. The specific value of \( c \) can be provided only when the following solvability condition is satisfied:

\[ (R'_0, M[R'_0]) \equiv \int_{-\infty}^{\infty} dX R'_0 M[R'_0] = 0, \]

(22)

where \( M[R'_0] = (1/g_1) \left( g_3 \partial_x^3 R' + g_1 \partial_x^2 R' + (g_1 g_3/g_2) \partial_x^2 R^3 \right). \) By solving equation (22), the general solution of propagation velocity \( c \) can be calculated as

\[ c = \frac{5g_3 g_1}{2g_2 g_4 - 3g_1 g_5}. \]

(23)

Hence, the corresponding kink-antikink solution is derived as

\[ \rho_j = \rho_c + \varepsilon \sqrt{\frac{g_1}{g_2}} \tanh \left( \frac{c}{\sqrt{2}} \left( X - c\rho \right) \right), \]

(24)

where \( \varepsilon^2 = ((a/a_0) - 1) \), and the amplitude of the density wave \( A \) is defined as \( \sqrt{\left( g_1 \varepsilon^2 C/g_3 \right)} \). According to the kink-antikink solution equation (24), one can get the coexisting phases which include the freely moving phase with low density \( (\rho_j = \rho_c - A) \) and the congested phase with high density \( (\rho_j = \rho_c + A) \).

Based on the linear and nonlinear theory analyses, numerical simulations are carried out with the comprehensive consideration of the anticipation effect of flux difference and the “backward-looking” effect.

4. Numerical Simulations

As a positive constant \( \gamma \) considering the “backward-looking” effect has been studied fully in some previous studies \([36, 40, 41]\), this paper will focus on the effects of parameters \( p \) and \( \tau \) in this part. Periodic boundary conditions are adopted by a step function as follows:

\[ \rho_j(0) = \rho_0 = 0.25, \quad j \in [1, N], \]

\[ \rho_j(1) = \rho_0 + \Delta \rho: \quad j = \frac{N}{2} - 1, \]

\[ \rho_j - \Delta \rho: \quad j = \frac{N}{2}, \]

where \( \Delta \rho = 0.01 \) is the initial perturbation and other parameters are set as \( N = 200 \), \( a = 1.2 \), and \( \rho_0 = \rho_c = 0.25 \).

4.1. The Effect of Anticipation Time \( \tau \). In order to analyze the anticipation time’s influence on the stability of the traffic system, the temporal and spatial evolution of density waves between time \( t = 10000 - 10200 \) s is depicted in Figure 4 for different anticipation time \( \tau \) with fixed parameters \( p = 0.1 \) and \( \gamma = 0.05 \). Since the given parameters in Figures 4(a)–4(c) do not satisfy the linear stable condition equation (14), kink-antikink density waves appear and the initial stable traffic flow evolves into nonuniform traffic flow with the added perturbation. To reveal the amplitude of the density wave more clearly, Figure 5 describes the density profile of each

\[ g_1 = \frac{1}{6} \rho_c \left( 7V'_b + V'_f + 3 \rho_p \varepsilon^2 (V'_b + V'_f)^3 \right), \]

\[ g_2 = \rho_c^2 \left( V'_b + V'_f \right), \]

\[ g_3 = \frac{1}{2} \rho_c^2 \left( 3V'_b - V'_f - 2 \rho (V'_b + V'_f) + 2 \rho_p \varepsilon^2 (V'_b + V'_f)^3 \right), \]

\[ g_4 = -\frac{1}{24} \rho_c^2 \left( 15V'_b - V'_f + 4 \rho_p \varepsilon^2 (V'_b + V'_f)^3 \right) - \frac{1}{3a_c^4} \left( a \rho \right)(a \rho - 1)(V'_b + V'_f) \left( 3 \rho_p \varepsilon^2 (V'_b)^3 + V'_f \right) \]

\[ -\frac{1}{3a_c^4} \left( a \rho \right)(a \rho - 1)(V'_b + V'_f) \left( 3 \rho_p \varepsilon^2 (V'_b)^3 + V'_f \right), \]

\[ g_5 = -\frac{1}{12} \rho_c^2 \left( 3V'_b - V'_f \right) - \frac{1}{3a_c^4} \left( a \rho \right)(a \rho - 1)(V'_b + V'_f) \left( V''_b + V''_f \right). \]
Figure 4: The evolution of the traffic densities for different values of parameter $\tau$: (a) $\tau = 1.0$; (b) $\tau = 1.5$; (c) $\tau = 2.0$; (d) $\tau = 2.5$.

Figure 5: The density profile at $t = 10200$ with different values of $\tau$: (a) $\tau = 1.0$; (b) $\tau = 1.5$; (c) $\tau = 2.0$; (d) $\tau = 2.5$. 
Figure 6: The evolution of the traffic densities for different values of parameter \( p \): (a) \( p = 0 \); (b) \( p = 0.05 \); (c) \( p = 0.1 \); (d) \( p = 0.15 \).

Figure 7: Continued.
subfigure shown in Figure 4 at $t=10200$. Moreover, Figures 5(a)–5(c) reveal that the amplitudes of density waves corresponding to Figures 4(a)–4(c) weaken as the anticipation time $\tau$ increases. When $\tau \approx 2.5$, i.e., the parameters in Figure 4(d) satisfy the linear stability condition, the added perturbation will disappear, and the traffic flow will eventually return to a uniform state. Figure 5(d) can get results similar to Figure 4(d). This phenomenon indicates that driver’s anticipation effect of flux difference could significantly promote the stabilization of traffic flow.

4.2. The Effect of Sensitivity Coefficient $p$. To expose the sensitivity coefficient’s influence on the stability of the traffic system, the temporal and spatial evolution of density waves between time $t = 10000 – 10200$ s is exhibited in Figure 6 for different values of $p$ with fixed $\gamma = 0.05$ and $\tau = 2$. When $p = 0$, the FADBDE model degrades into Ge and Cheng’s model [39]. Note that the parameters in Figures 6(a)–6(c) do not satisfy the linear stable condition equation (14). Therefore, the stop-go traffic density waves occur with the small added perturbation in Figures 6(a)–6(c). Meanwhile, the amplitude of the density wave decreases gradually as the value of $p$ increases. Especially in Figure 6(d), the congested state entirely disappears when $p \approx 0.15$. The relationship between Figures 6 and 7 is similar to the relationship between Figures 4 and 5. Figure 7 describes the density profile of each subfigure shown in Figure 6 at $t=10200$. As parameter $p$ increases, the traffic flow becomes more gentle and stable. This indicates that the sensitivity coefficient $p$ contributes to suppressing the traffic congestion.

Figure 8(a) and 8(b) describe the flux against the density with different parameters corresponding to the traffic flow in Figures 4 and 6, respectively. In Figure 8, the amplitude of the hysteresis loops reduced with the increasing value of $\tau$ or $p$. When the corresponding parameter $\tau \geq 2.5$ or $p \geq 0.15$, the hysteresis loop turns into a stable point, which means that the traffic flow eventually becomes stable even with the
added small perturbation. To sum up, as the anticipation time $\tau$ or the sensitivity coefficient $p$ increases, the stability of the system can be enhanced by incorporating the driver's anticipation effect and the "backward-looking" effect.

5. Conclusion

In this paper, a novel lattice hydrodynamic model with the anticipation effect of flux difference and the "backward-looking" effect is constructed to explore the traffic driving behavior for a single-lane traffic system. The stability conditions and the solution of the mKdV equation are deduced, respectively, based on the linear stability analysis and nonlinear theory. The results of the theoretical analysis display that the increasing values of both the driver's anticipation time and the sensitivity coefficient of the anticipation flux difference can contribute to enlarging the stable region. Numerical simulations reveal that the FADBEm model is effective in stabilizing traffic flow even with perturbations. By considering the anticipation effect of flux difference and the "backward-looking" effect simultaneously, the traffic congestion can be alleviated effectively, which is in accordance with the conclusions of the theoretical analysis. Future work will focus on solving the traffic congestion on multilane road problem, the automation for longitudinal driving, and the vehicular emergency warning system.

Data Availability

The data used to support the findings of this study are available from the corresponding author upon request.

Conflicts of Interest

The authors declare that there are no conflicts of interest regarding the publication of this paper.

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