Dependence of lattice strain relaxation, absorbance, and sheet resistance on thickness in textured ZnO@B transparent conductive oxide for thin-film solar cell applications

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Abstract
The interplay of surface texture, strain relaxation, absorbance, grain size, and sheet resistance in textured, boron-doped ZnO (ZnO@B), transparent conductive oxide (TCO) materials of different thicknesses used for thin film, solar cell applications is investigated. The residual strain induced by the lattice mismatch and the difference in the thermal expansion coefficient for thicker ZnO@B is relaxed, leading to an increased surface texture, stronger absorbance, larger grain size, and lower sheet resistance. These experimental results reveal the optical and material characteristics of the TCO layer, which could be useful for enhancing the performance of solar cells through an optimized TCO layer.

Introduction
Thin-film solar cells require a transparent conductive oxide (TCO) to allow light to reach the absorber layers and create the electrical current. Due to its superior characteristics, including a wide band gap, high dielectric constant, high exciton binding energy (60 meV), high thermal stability, high transparency, and high conduction, wurtzite ZnO is a very promising TCO material used for the front contact, barrier layer, and intermediate reflector in solar cells [1-9]. Low-pressure chemical vapor deposition (LPCVD) can be implemented to deposit such a transparent, textured, and highly conductive TCO [9]. For thin-film solar cell applications, the LPCVD-grown ZnO can possess an as-grown textured structure to enhance light scattering and to increase the optical path through the solar cell without any post-treatment. Because the structural, optical, and electrical characteristics of LPCVD-grown ZnO are sensitive to the growth temperature, pressure, and flow rate, TCO can be tuned according to the application.

ZnO films grown on a sapphire substrate undergo residual strain induced by the lattice mismatch and the difference in thermal expansion coefficients at the interface, leading to strain relaxation in the material. This strain relaxation influences the optical and electrical properties of the TCO layer, such as absorbance, sheet resistance, and grain size. Therefore, understanding the interplay of these factors is crucial for improving the performance of thin-film solar cells.
expansion coefficient [10]. Because strain can affect the electronic and optical properties of materials, the strain distribution in the films is an important subject to be investigated. The strain in ZnO films is accumulated during film growth and can be monitored by in situ optical reflectance measurement as a function of thickness, \( t \): compressive strain for \( t < 5.5 \) nm, released compressive strain due to generation of misfit dislocations for \( 5.5 \text{ nm} < t < 200 \) nm, tensile strain due to thermal stress for \( 200 \text{ nm} < t < 500 \) nm, and residual tensile strain relaxed by microcrack formation for \( t > 500 \) nm [10].

The variation in the physical properties of nanostructures drastically influences the optoelectronic properties of ZnO [11-13]. X-ray-excited optical luminescence of ZnO nanoneedles shows a sharp band gap emission and a broad red emission related to surface defects, while that of ZnO microcrystallites has a strong green emission due to defect states in the core [11]. A blue-shifted absorption edge and photoluminescence caused by quantum confinement as well as a higher photovoltaic and sensor performance due to a larger surface area have been demonstrated in ZnO nanocrystals [14-16]. In addition, the microstructure, optical properties, and strain of thickness-dependent ZnO thin film grown by atomic layer deposition have been reported [17]. The thicker ZnO thin films show a larger crystalline grain, a smaller lattice strain, a higher Zn/O ratio, and better crystal quality. Furthermore, doping impurities, such as B, Al, and Ga, can improve the electrical transport properties of ZnO [18].

There are two benefits to using a B-doped ZnO (ZnO@B) film as the TCO layer. First, B has the smallest ionic radius among the three dopants (B\(^{3+}\): 0.23 Å, Al\(^{3+}\): 0.54 Å, Ga\(^{3+}\): 0.62 Å), which results in better transparency. Second, a textured surface can be easily achieved for a ZnO@B film grown by LPCVD, which can enhance light scattering and increase the optical path through the solar cell without any post-treatment. However, many important issues regarding optical and material characteristics in textured, ZnO@B TCO grown by LPCVD for thin-film solar cell applications is yet to be explored.

This study reports the lattice strain relaxation, absorbance, and sheet resistance of textured ZnO TCO@B for solar cell applications. A thicker ZnO@B film enhances the strain relaxation, resulting in an increased surface texture, stronger absorbance, larger grain size, and lower sheet resistance. The optimization of the TCO layer could be useful for enhancing the performance of solar cells.

**Results and Discussion**

**Structural characterization**

Four ZnO@B samples with 20-, 40-, 60-, and 70-minute growth times were prepared (here within named, c-20, c-40, c-60, and c-70, respectively). Figure 1 shows the atomic force microscopy (AFM) images (5 × 5 μm) of the four ZnO@B samples. The surface roughness is 16.603, 26.756, 51.531 and 56.233 nm for c-20, c-40, c-60, and c-70, respectively. Each sample is composed of small grains. The grain size of the triangular features is estimated to be 676, 1520, 2706, and 4220 nm\(^2\) for c-20, c-40, c-60, and c-70, respectively. As the growth time (thickness) increases, the surface becomes more textured and the grain size larger. The film is composed of vertically stacked monocrystalline grains that appear as pyramids at the surface. The apparent texture structure in the sample c-70 can effectively scatter light to enhance the light coupling.

Figure 1a–d shows scanning electron microscope (SEM) images of the c-20, c-40, c-60, and c-70 ZnO@B samples, respectively. Figure 2a–d shows the panchromatic cathodoluminescence (CL) images of the corresponding SEM regions using an 11 kV excitation electron voltage. In the SEM images, the c-20 sample is composed of small grains with no preferential orientation. Increasing thickness leads to an apparent textured structure and larger grain size. The CL image of the c-70 sample shows a high contrast and bright image, while that of the c-20 shows a low contrast and dark one. This shows that the apparent texture structure enhances light scattering.

Figure 3a–d shows the cross-sectional SEM images for the c-20, c-40, c-60, and c-70 ZnO@B samples, respectively. From the cross-sectional SEM images of each sample, the thicknesses of the c-20, c-40, c-60, and c-70 samples are estimated to be 271.9, 543.7, 1,022, and 1388 nm, respectively.
Optical properties

Figure 4 shows the CL spectra of the four samples at room temperature (RT). As the thickness increases, the CL emission peak is slightly red-shifted. An emission peak around 378 nm (3.28 eV) is related to a band-to-band transition, while that in the 382–386 nm (3.21–3.24 eV) spectral range is due to transitions between band-tail states of ZnO [17]. Hence, as the thickness increases, the dominant red-shifted CL emission moves from the band-to-band transition to the transition between the band-tail states. Due to a larger grain size in the thicker c-70 sample, the red-shifted CL emission is consistent with the weaker quantum size effect.

Figure 5 shows the absorbance squared, $\alpha^2(E)$, of the four samples. Compared to the c-20 sample, the other three samples show enhanced absorbance in the 300–350 nm spectral range. As the thickness increases, the sharper, red-shifted absorption edge in the 360–380 nm spectral range is consistent with the
more uniform distribution of grain size and weaker quantum size effect. The band gap energy, $E_g$, of a semiconductor can be estimated by extrapolating the linear portion of the absorbance square to zero. As shown in Table 1, as the thickness increases, the red-shifted $E_g$ is consistent with the weaker quantum size effect. The lower $E_g$ of ZnO@B samples compared to that of Al-doped ZnO (AZO or ZnO@Al) samples is beneficial for photon absorption and contributes to the quantum efficiency of solar cells [19].

### Strain relaxation

To study the effect of thickness on strain relaxation, Raman measurements were taken to estimate the in-plane strain of the four samples. Figure 6 shows the Raman scattering spectra for the four samples. The spectra display $A_1$(TO) and $A_1$(LO) modes for ZnO. The dotted lines at 379 cm$^{-1}$ and 574 cm$^{-1}$ show the strain-free $A_1$(TO) and $A_1$(LO) modes, respectively, for ZnO [20]. The in-plane strain in the $x$-direction, $\varepsilon_{xx}$, and in the $y$-direction, $\varepsilon_{yy}$, of ZnO films can be determined by the frequency shift, $\Delta \omega = \omega - \omega_0$ [21], as:

$$\varepsilon_{xx} = \varepsilon_{yy} = \frac{\Delta \omega (A_1(\text{TO}))}{b - a} \frac{C_{13}}{C_{33}} \frac{C_{13}}{2C_{13}}$$

where $a = -774$ cm$^{-1}$ and $b = -375$ cm$^{-1}$ are the deformation potential constants of the $A_1$(TO) mode [22]. The elastic stiffness constants, $C_{33}$ and $C_{13}$, are 216 and 104 GPa, respectively [1]. The six-fold symmetry of the hexagonal polar $c$-ZnO dictates an isotropic in-plane strain in the basal plane, i.e., $\varepsilon_{xx} = \varepsilon_{yy} = \varepsilon$. The frequency shift, $\Delta \omega$, for the $A_1$(TO) mode are shown in Table 1. The in-plane strain $\varepsilon$ can be deduced from Equation 1.

The in-plane tensile strain, $\varepsilon$, of the $c$-20, $c$-40, $c$-60, and $c$-70 samples is 0.0069, 0.0682, 0.1292, and 0.0069$, respectively. Figure 7 shows the evolution of in-plane strain for the four samples as a function of thickness. As the thickness increases, the in-plane tensile strain becomes larger and then relaxed. A similar trend, attributed to strain relaxation by microcrack forma-

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**Table 1:** Thickness (nm), grain size (nm$^2$), roughness (nm), $E_g$ (eV), $\Delta \omega (A_1(\text{TO}))$ (cm$^{-1}$), $\varepsilon$ (%), and sheet resistance ($R_{\text{sheet}}$) (ohm/□) for the four ZnO@B samples.

| Sample | Thickness (nm) | Grain size (nm$^2$) | Roughness (nm) | $E_g$ (eV) | $\Delta \omega (A_1(\text{TO}))$ (cm$^{-1}$) | $\varepsilon$ (%) | $R_{\text{sheet}}$ (ohm/□) |
|--------|----------------|---------------------|----------------|-----------|---------------------------------------------|------------------|----------------------|
| c-20   | 272            | 696                 | 16.603         | 3.32      | $-0.082$                                    | 0.0069           | 846                  |
| c-40   | 544            | 1,220               | 26.756         | 3.31      | $-0.809$                                    | 0.0682           | 191                  |
| c-60   | 1,022          | 2,706               | 51.531         | 3.28      | $-1.534$                                    | 0.1292           | 96                   |
| c-70   | 1,388          | 4,220               | 56.233         | 3.26      | $-0.082$                                    | 0.0069           | 76                   |
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Conclusion
In summary, a thicker, ZnO@B TCO layer enhances the strain relaxation and hence shows an increased surface texture, stronger absorbance, less transparency, larger grain size, and lower sheet resistance. For solar cell applications, the TCO layer should be highly transparent and less resistive. Although the smaller sheet resistance and increased surface texture in thicker ZnO@B samples are beneficial to device performance, photons in the UV spectral range are strongly absorbed in such a TCO layer and could limit the quantum efficiency of the thin-film solar cell. For the application of ZnO@B TCO in solar cells, the TCO layer must be grown on a solar cell and its optimal thickness therefore could be chosen to match the overall material and structure of the cell for enhancing the performance of solar cells. Nevertheless, this is an interesting and important subject which will be left for future investigation.

Experimental Synthesis of ZnO@B thin films by LPCVD
The polar c-plane ZnO samples were grown on a sapphire substrate at 170 °C at a pressure of 0.6 Torr in a LPCVD reactor. DEZ and H₂O were used as the precursors for Zn and O, respectively. ZnO was doped with B to improve the electrical transport properties of ZnO thin films [9,18]. B₂H₆ was used as a doping gas. Four ZnO@B samples with 20-, 40-, 60-, and 70-minute growth times, corresponding to 272, 544, 1022, and 1388 nm in thickness, were prepared (samples c-20, c-40, c-60, and c-70, respectively). For the four samples, the B₂H₆, DEZ, and H₂O flow rates were 1, 500, and 550 sccm, respectively.

Characterization
The surface morphology was revealed by atomic force microscopy (Park Systems, XE-70) performed in noncontact mode using a silicon tip with a curvature of less than 10 nm. The scanning electron microscope and cathodoluminescence results were acquired by the use of a Gatan monoCL3 spectrometer in a JEOL JSM 7000F SEM system. The absorption spectra were acquired with a U-3900 spectrophotometer (model 2J2-0015) at room temperature. The Raman spectra were recorded in the backscattering configuration using a Jobin Yvon-Horiba micro-Raman system (model T64000) with a 532 nm laser.

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References
1. Morkoç, H.; Özgür, Ü. Zinc Oxide: Fundamentals, Materials and Device Technology; Wiley-VCH: Weinheim, Germany, 2007. doi:10.1002/9783527623945
2. Özgür, Ü.; Alivov, Ya. I.; Liu, C.; Teke, A.; Reschkeikov, M. A.; Doğan, S.; Avrutin, V.; Cho, S. J.; Moroć, H. J. Appl. Phys. 2005, 98, 041301. doi:10.1063/1.1992666

3. Lai, C. M.; Huang, Y. E.; Kou, K. Y.; Chen, C. H.; Tu, L. W.; Feng, S. W. Appl. Phys. Lett. 2015, 107, 022110. doi:10.1063/1.4926678

4. Söderström, T.; Domine, D.; Feltrin, A.; Despeisse, M.; Meillaud, F.; Bugnon, G.; Boccard, M.; Cuony, P.; Haug, F. J.; Fajé, S.; Nicolay, S.; Ballif, C. Proc. SPIE 2010, 7603, 76030B. doi:10.1117/12.843511

5. Chen, X. L.; Liu, J. M.; Ni, J.; Zhao, Y.; Zhang, X. D. Appl. Phys. Lett. 2011, 98, 041301. doi:10.1063/1.1992666

6. Lai, C. M.; Huang, Y. E.; Kou, K. Y.; Chen, C. H.; Tu, L. W.; Feng, S. W. Appl. Phys. Lett. 2015, 107, 022110. doi:10.1063/1.4926678

7. Bikowski, A.; Ellmer, K. J. Mater. Res. 2012, 27, 2249–2256. doi:10.1557/jmr.2012.113

8. Slassi, A.; Ziel, Y.; Zarhri, Z.; Abdellaoui, M.; Fakhim Lamrani, A. Phys. Scr. 2015, 90, 085801. doi:10.1088/0031-8949/90/8/085801

9. Nikitina, G.; Glazunov, N.; Fedyanin, Y.; Al'tshuler, E. Phys. Scr. 2015, 90, 085801. doi:10.1088/0031-8949/90/8/085801

10. Bikowski, A.; Ellmer, K. J. Mater. Res. 2012, 27, 2249–2256. doi:10.1557/jmr.2012.113

11. Arno, H.; Morkoç, H.; Mierendorf, W.; Macek, E.; Wrede, B. J. Appl. Phys. 2007, 91, 231904. doi:10.1063/1.2813021

12. Amelio, L.; Heilg, F.; Brunet, S.; Sammenaiken, R.; Regier, T.; Blyth, R. I. R.; Zuin, L.; Sankari, R.; Vogt, J.; Sham, T. K. ChemPhysChem 2010, 11, 3625–3631. doi:10.1002/cphc.201000730

13. Dutta, S.; Chattopadhyay, S.; Sarkar, A.; Chakrabarti, M.; Sanyal, D.; Jana, D. Prog. Mater. Sci. 2009, 54, 89–136. doi:10.1016/j.pmatsci.2008.07.002

14. Li, D.; Leung, Y. H.; Durdzic, A. B.; Liu, Z. T.; Xie, M. H.; Shi, S. L.; Xu, S. J.; Chan, W. K. Appl. Phys. Lett. 2004, 85, 1601–1603. doi:10.1063/1.1786375

15. Vanheusden, K.; Seager, C. H.; Warren, W. L.; Tallant, D. R.; Voigt, J. A. Phys. Appl. Lett. 1996, 68, 403–405. doi:10.1063/1.116699

16. Sun, C. Q. Prog. Solid State Chem. 2007, 35, 1–159. doi:10.1016/j.progsolidstchem.2006.03.001

17. Chaya, A. A.; Viter, R.; Bechelany, M.; Alute, Z.; Ertz, D.; Zalesskaya, A.; Kovalevskis, K.; Roussac, V.; Smytyna, V.; Miele, P. Beilstein J. Nanotechnol. 2013, 4, 690. doi:10.3762/bjnano.4.78

18. Wang, G.; Li, Z.; Li, S.; Li, M.; Shi, C.; Liao, J.; Chen, C. Ceram. Int. 2016, 42, 2813–2820. doi:10.1016/j.ceramint.2015.11.014

19. Xiao, S.; Zhou, J.; Huang, S.; Xiao, P.; Gu, X.; Yan, D.; Xu, S. J. Phys. D. Appl. Phys. 2015, 48, 305105. doi:10.1088/0022-3727/48/30/305105

20. Alim, K. A.; Fonoberov, V. A.; Shamsa, M.; Balandin, A. A. J. Appl. Phys. 2005, 97, 124313. doi:10.1063/1.1944222

21. Gruber, T.; Prinz, G. M.; Kirchner, C.; Kling, R.; Reuss, F.; Limmer, W.; Waag, A. J. Appl. Phys. 2004, 96, 289–293. doi:10.1063/1.1755433

22. Calsen, G.; Reparaz, J. S.; Wagner, M. R.; Kirsch, R.; Nenstiel, C.; Hoffmann, A.; Phillips, M. R. Appl. Phys. Lett. 2011, 98, 061906. doi:10.1063/1.3554434