Mapping images into ordinal networks

Arthur A. B. Pessa* and Haroldo V. Ribeiro†

Departamento de Física, Universidade Estadual de Maringá – Maringá, PR 87020-900, Brazil

(Dated: July 8, 2020)

An increasing abstraction has marked some recent investigations in network science. Examples include the development of algorithms that map time series data into networks whose vertices and edges can have different interpretations, beyond the classical idea of parts and interactions of a complex system. These approaches have proven useful for dealing with the growing complexity and volume of diverse data sets. However, the use of such algorithms is mostly limited to one-dimension data, and there has been little effort towards extending these methods to higher-dimensional data such as images. Here we propose a generalization for the ordinal network algorithm for mapping images into networks. We investigate the emergence of connectivity constraints inherited from the symbolization process used for defining the network nodes and links, which in turn allows us to derive the exact structure of ordinal networks obtained from random images. We illustrate the use of this new algorithm in a series of applications involving randomization of periodic ornaments, images generated by two-dimensional fractional Brownian motion and the Ising model, and a data set of natural textures. These examples show that measures obtained from ordinal networks such as average shortest path and global node entropy extract important image properties related to roughness and symmetry.

I. INTRODUCTION

In the last two decades, network science has established itself as a vibrant and successful field of research [1]. The fact that diverse complex systems are accurately described as sets of vertices and edges [2, 3] combined with advancements in data acquisition and processing have led to the widespread application of networks to an immense variety of systems from biological, sociological and technological origins. In addition to these more established works, recent developments in network science have originated new and more abstract forms to define complex networks. Vertices and edges of these more abstract networks usually do not represent constituents and interactions of a system; instead, algorithms mapping objects into networks ascribe different meanings to these basic network components [4, 5].

Algorithms designed to map time series into networks are a particularly important class of such networks [5] which have been used to characterize time series of chaotic and stochastic nature obtained from simulations and experimental data. These approaches are collectively known as time series networks and the most prominent examples include visibility graphs [6], recurrence networks [7] and ordinal networks [8]. Visibility graphs and its variants [6, 9–12] map each time series observation into a network vertex, and any two vertices are connected if their respective values in the time series satisfy a visibility condition [6]. Recurrence networks, on its turn, are a reinterpretation of recurrence plots [13], a graphic tool developed in the context of nonlinear time series analysis [14, 15]. Vertices of recurrence networks represent a set of states obtained from small segments of time series, and edges connect vertices (pairs of states) that are similar according to a distance metric [7]. Ordinal networks were proposed more recently but also originate from the study of nonlinear time series, specifically from a successful symbolization approach due to Bandt and Shiha [16] and Bandt and Pompe [17]. Nodes of ordinal networks represent ordering patterns (or permutation symbols) associated with time series slices, and links are drawn based on the succession of these symbols in the time series [8, 18–20].

While the use of algorithms mapping time series into networks is increasingly gaining popularity among the scientific community (see Zou et al. [5] for a recent review), few efforts have been dedicated to extending these ideas to higher-dimensional data such as images. Notable exceptions include the works of Xiao et al. [21] on the degree distributions of row-column visibility graphs obtained from fractal landscapes, and Lacasa and Iaco-vacci [22, 23] that analyzed visibility networks mapped from bidimensional random and chaotic data and tackled problems of image processing and classification. As processes for extracting meaningful information from images are ubiquitous across science, extending and proposing approaches for mapping images into complex networks have great potential to contribute with new image quantifiers derived from well-established network metrics.

Here we present an extension of the ordinal network framework allowing the representation of images as complex networks. We describe intrinsic connectivity constraints of ordinal networks inherited from the symbolization process and determine the exact form of ordinal networks mapped from completely noisy images. By mapping images obtained from periodic ornaments, two-dimensional fractional Brownian motion, and the Ising model into ordinal networks, we illustrate the use of this new approach to identify and describe these systems with network-related metrics. We further apply our method
to characterize a data set of real-world images, where we show that ordinal networks are capable of distinguishing different types of textures and identifying image symmetries.

The rest of this paper is organized as follows. In Section II, we briefly revisit the ordinal network framework before generalizing it to two-dimensional data. Next, we investigate connectivity properties of ordinal networks (Section III A), the exact form of ordinal networks obtained from random data (Section III B), and noisy-periodic ornaments (Section III C). Applications involving fractal and Ising surfaces are presented in Sections III D and III E, and real-world images (Brodatz textures) are investigated in Section III F. Finally, Section IV concludes our work.

II. METHODS

The representation of time series as ordinal networks has direct relations to permutation entropy, a successful time series complexity measure [17]. Specifically, ordinal networks use the same approach introduced by Bandt and Shiha [16] and Bandt and Pompe [17] to partition a time series in small segments and associate a permutation symbol (or an ordinal pattern) with each part according to the relative amplitude of the time series values [19]. Our extension of the ordinal network algorithm for two-dimensional data is inspired by a generalization of permutation entropy to image data proposed in Ref. [24] that proved useful for investigating liquid crystals [25, 26] and art paintings [27].

By following Ref. [24], we start by considering a two-dimensional array \( \{ y_{ij} \}_{i=1,...,N_x} \times_{j=1,...,N_y} \) of size \( N_x \times N_y \), where the elements \( y_{ij} \) may represent pixels of an image. Next, we divide this array into sliding partitions of size \( d_x \) by \( d_y \) defined as

\[
\begin{align*}
w_s^t &= \left( \begin{array}{cccc}
  y_{s} & y_{s+1} & \ldots & y_{s+(d_y-1)} \\
  y_{s+1} & y_{s+2} & \ldots & y_{s+(d_y-1)} \\
  \vdots & \vdots & \ddots & \vdots \\
  y_{s+(d_x-1)} & y_{s+(d_x-1)+1} & \ldots & y_{s+(d_x-1)+(d_y-1)}
\end{array} \right),
\end{align*}
\]

(1)

where the indices \( s = 1, \ldots, N_x \) and \( t = 1, \ldots, N_y \), with \( N_x = N_x - d_x + 1 \) and \( N_y = N_y - d_y + 1 \), cover all possible sliding partitions. The values of \( d_x \) and \( d_y \) are the two parameters of the approach and represent the horizontal and vertical embedding dimensions [24]. We then flatten these two-dimensional partitions line by line as

\[
\begin{align*}
w_s^t &= (y_s, y_{s+1}, \ldots, y_{s+(d_y-1)}, \\
y_{s+1}, y_{s+1+1}, \ldots, y_{s+1+(d_y-1)}, \ldots, \\
y_{s+(d_x-1)}, y_{s+(d_x-1)+1}, \ldots, y_{s+(d_x-1)+(d_y-1)}).
\end{align*}
\]

(2)

to investigate the ordering of its elements. Because this procedure does not depend on the partition location (that is, \( s \) and \( t \)), we can simplify the notation and rewrite the flattened partition as

\[
w = (\tilde{y}_0, \tilde{y}_1, \ldots, \tilde{y}_{d_xd_y-2}, \tilde{y}_{d_xd_y-1}),
\]

(3)

where \( \tilde{y}_0 = y_0^t, \tilde{y}_1 = y_1^t+1, \) and so on.

Under this notation, the symbolization procedure consists in evaluating the permutation \( \Pi = (r_0, r_1, \ldots, r_{d_xd_y-2}, r_{d_xd_y-1}) \) of the index numbers \( (0, 1, \ldots, d_xd_y-2, d_xd_y-1) \) that sorts the elements of the flattened partition in ascending order, that is, the index numbers resulting in \( \tilde{y}_r \leq \tilde{y}_r \leq \cdots \leq \tilde{y}_{d_xd_y-2} \leq \tilde{y}_{d_xd_y-1} \). In case of draws, we maintain the occurrence order of the elements in \( w \), that is, \( r_{j-1} < r_j \) if \( \tilde{y}_{r_{j-1}} = \tilde{y}_{r_j} \) for \( j = 1, \ldots, d_xd_y - 1 \) [28, 29]. To illustrate this procedure, suppose we have \( d_x = d_y = 2 \) and the partition matrix \( w = \left[ \begin{array}{cc} 4 & 5 \\ 2 & 4 \end{array} \right] \). The corresponding flattened array is \( w = (4, 5, 2, 2) \), and so \( \tilde{y}_0 = 4, \tilde{y}_1 = 5, \tilde{y}_2 = 2 \) and \( \tilde{y}_3 = 2 \). Because \( \tilde{y}_2 \leq \tilde{y}_3 \leq \tilde{y}_0 \leq \tilde{y}_1 \), the permutation \( \Pi = (2, 3, 0, 1) \) is the one that sorts the elements of the partition \( w \).

After carrying out the symbolization procedure over the entire data array, we construct another array \( \{ \pi_i \}_{i=1,...,n_y} \) containing the permutation symbols associated with each sliding partition \( w_i^t \). By using this new array, we can calculate the relative frequency \( \rho_i(\Pi_i) \) of each possible permutation \( \Pi_i \) defined as

\[
\rho_i(\Pi_i) = \frac{\text{number of partitions of type } \Pi_i \text{ in } \{ \pi_i \}}{n_xn_y},
\]

(4)

where \( i = 1, \ldots, (d_xd_y)! \) and \( (d_xd_y)! \) is the total number of possible permutations that can occur in the original data array. Having these relative frequencies, we construct the probability distribution \( P = \{ \rho_i(\Pi_i) \}_{i=1,...,(d_xd_y)!} \) of ordinal patterns and estimate the two-dimensional version of the permutation entropy [24]

\[
H = - \sum_{i=1}^{(d_xd_y)!} \rho_i(\Pi_i) \log \rho_i(\Pi_i),
\]

(5)

where \( \log(\ldots) \) stands for the base-2 logarithm. It is worth noticing that the embedding dimensions \( d_x \) and \( d_y \) must satisfy the condition \( (d_xd_y)! \ll N_xN_y \) in order to obtain a reliable estimate of the probability distribution \( P = \{ \rho_i(\Pi_i) \}_{i=1,...,(d_xd_y)!} \) [17, 24]. It is worth mentioning that this generalized version of the permutation entropy recovers the one-dimensional case (time series data or \( N_y = 1 \)) by setting \( d_y = 1 \) and properly choosing \( d_x \).

To generalize the concept of ordinal networks to two-dimensional data, we use the symbolic array \( \{ \pi_i \}_{i=1,...,n_y} \) obtained from the previous discussion. As in the one-dimensional case, we consider each unique permutation symbol \( \Pi_i \), occurring in \( \{ \pi_i \} \) as a node of the corresponding ordinal network. Next, we draw directed edges between these nodes according to the first-neighbor transitions occurring in \( \{ \pi_i \}_{i=1,...,n_y} \) that is, we directly-connect the permutation symbols involved in all
horizontal and vertical successions among ordinal patterns in the symbolic array ($\pi^t_s \rightarrow \pi^{t+1}_s$ and $\pi^t_s \rightarrow \pi^{t+1}_s$, with $s = 1, \ldots, n_x - 1$ and $t = 1, \ldots, n_y - 1$). The directed link between a pair of permutation symbols $\Pi_i$ and $\Pi_j$ is weighted by the total number of occurrences of this particular transition in the symbolic array. Thus, we can write the elements of the weighted-adjacency matrix representing the ordinal network as

$$ p_{i,j} = \frac{\text{total of transitions } \Pi_i \rightarrow \Pi_j \text{ in } \{\pi^t_s\}_{s=1,\ldots,n_x}^{t=1,\ldots,n_y}}{2n_x n_y - n_x - n_y}, \quad (6) $$

where $i, j = 1, \ldots, (d_x d_y)!$ and the denominator represents the total number of horizontal and vertical permutation successions in $\{\pi^t_s\}_{s=1,\ldots,n_x}^{t=1,\ldots,n_y}$. Figure 1 illustrates the procedure for creating an ordinal network from a simple two-dimensional array of size $N_x = 3$ and $N_y = 4$.

In addition to more usual network metrics, the probabilistic aspects of nodes and edges in ordinal networks can also be quantified by entropy measures at the node level or for the whole network. Given an ordinal network vertex $i$ (associated with a permutation $\Pi_i$), the local node entropy $H_i$ is defined for this vertex as

$$ h_i = -\sum_{j \in \mathcal{O}_i} p'_{i,j} \log p'_{i,j}, \quad (7) $$

where $p'_{i,j} = p_{i,j}/\sum_{k \in \mathcal{O}_i} p_{i,k}$ represents the renormalized transition probability of transitioning from node $i$ to node $j$ (associated with the permutation $\Pi_j$), and $\mathcal{O}_i$ is the outgoing neighborhood of node $i$ (set of all edges leaving node $i$). The local node entropy $h_i$ quantifies the degree of determinism related to permutation transitions at the node level. We have $h_i = 0$ (deterministic case) when only one edge leaves node $i$, whereas $h_i$ is maximum if all edges leaving $i$ have the same weight (equiprobable case). At the network level, we can define the global node entropy as

$$ H_{\text{GN}} = \sum_{i=1}^{(d_x d_y)!} p'_i h_i, \quad (8) $$

where $p'_i = \sum_{j \in \mathcal{I}_i} p_{i,j}$ corresponds to the probability of transitioning to node $i$ from its incoming neighborhood $\mathcal{I}_i$ (in-strength of node $i$). If the original data array is large enough ($|d_x d_y|! \gg N_x N_y$), $p'_i$ converges to the probability of occurrence of permutation $\pi_i$, and $H_{\text{GN}}$ corresponds to a weighted average of local node determinism throughout the network.

### III. RESULTS

#### A. Connectivity constraints

Having described our extension of the ordinal network algorithm for two-dimensional data, we start by investigating how properties of the symbolic array of permutations affect the connectivity of the resulting network. Similarly to what has been found for time series \cite{19}, ordinal networks mapped from two-dimensional arrays also present restrictions that prohibit the existence of several edges among permutation symbols. These constraints result from the symbolization approach and the procedure used to partition the array of data; they emerge even when data is completely random.

To better illustrate this finding, we consider a data array with $\{y_{ij}\}_{i=1,\ldots,N_x}^{j=1,\ldots,N_y}$ and embedding dimensions $d_x = d_y = 2$. Let us now suppose that the first partition matrix $w_1^1 = (y_1^1, y_2^1)$ is characterized by the permutation $\pi_1^1 = (2, 1, 3, 0)$ and that the next horizontally-adjacent partition is $w_1^2 = (y_1^2, y_2^2)$, $y_2^2 = y_2^1$. We note that both partitions share the array elements $y_1^2$ and $y_2^2$. Thus, because $w_1^1$ is described by $\pi_1^1 = (2, 1, 3, 0)$ (that is, $y_2^2 < y_1^1 < y_2^2 < y_1^1$), the permutation $\pi_1^2$ associated with $w_1^2$ must respect the inequality $y_1^1 < y_2^2$ imposed by $\pi_1^1$. Out of the $(d_x d_y)! = 24$ possible permutations, the previous constraint restricts $\pi_1^2$ to one among 12 permutations in which the index number 0 (corresponding to the position of $y_1^1$ in $w_1^2$) precedes the number 2 (corresponding to the position of $y_2^2$). Vertically-adjacent partitions also present similar constraints. For instance, $w_1^3 = (y_1^2, y_3^2)$ shares the array elements $y_2^2$ and $y_3^2$ with $w_1^1$, and the index number 0 (position of $y_1^1$ in $w_1^2$) has to precede the ordinal number 1 (index number of $y_2^2$ in $w_1^2$) in $\pi_1^3$ because of the inequality $y_2^2 < y_3^2$ expressed by $\pi_1^1$. Thus, $\pi_1^3$ is also constrained by $\pi_1^1$ to be one among 12 ordinal patterns. These restrictions hold independently of the particular permutation corresponding to $\pi_1^1$, and for each permutation, there are only 12 others that can immediately follow it (horizontally or vertically) when $d_x = d_y = 2$.

The same reasoning applies to permutation successions involving arbitrary values of the embedding dimensions $d_x$ and $d_y$. However, it is important to notice that the number of constraints in successions among permutations increases as adjacent partitions share a larger number of array elements. For instance, horizontally-adjacent partitions share two array elements while vertically-adjacent partitions share three array elements when $d_x = 3$ and $d_y = 2$. For these embedding dimensions, out of all $(3 \times 2)! = 720$ possible permutations, once the permutation $\pi_1^1$ is set, there are only 30 allowed ordinal patterns for its horizontal-neighbor permutation $\pi_1^{21}$ and 120 allowed ordinal patterns for its vertically-adjacent permutation $\pi_1^{31}$.

Once we know the set of allowed horizontal and vertical neighbors for a given permutation, the maximum number of outgoing edges in the vertex associated with this particular permutation is the intersection of these two sets. Interestingly, we have found that the maximum number of outgoing edges depends on the ordinal pattern itself. For instance, the ordinal pattern $\Pi = (0, 1, 2, 3)$ can have at most 16 outgoing edges while $\Pi = (0, 1, 3, 2)$ can have...
Arrays of data have all the same weight. Self-loops [as in the permutation Π of data array as an ordinal network. In this example, all permutation successions occur only once and so the network edges transitions (or vertical and horizontal successions) among ordinal patterns occurring in the symbolic array. (e) Representation sequences (or permutation patterns) \( s^1 \) and \( s^2 \) up to 20 outgoing connections in an ordinal network.

By investigating the maximum number of allowed edges for every permutation node, we can find the maximum number of edges for the whole ordinal network. For embedding dimensions \( d_x = d_y = 2 \), we find that the ordinal network resulting from an arbitrary data array can have up to 24 vertices linked by 416 edges. As in the one-dimensional case [19], the number of nodes and edges in ordinal networks increases dramatically with the embedding dimensions. For instance, an ordinal network can be formed by 720 nodes and 104,184 edges when \( d_x = 3 \) and \( d_y = 2 \). This imposes important limitations in the investigation of small data samples since the condition \( (d_x, d_y)! \ll N_x N_y \) is required to have a reliable estimate for the transition probabilities (edge weights). Most empirical investigations are therefore limited to values of embedding dimensions equal to two or three.

### B. Random ordinal networks

The constraints discussed in the previous section already indicate that ordinal networks emerging from completely random arrays are not random graphs. As it has been recently uncovered for time series [19], ordinal networks resulting from random data (called random ordinal networks) have a non-trivial network structure. We now investigate how these findings generalize for two-dimensional random data. To do so, we consider an array of data \( \{y^j_i \}_{i=1, \ldots, N_x} \) sampled from a continuous probability distribution and embedding dimensions \( d_x = d_y = 2 \). We then extract the first two horizontally-adjacent partitions \( w^1_i = \left( y^1_i, y^2_i \right) \) and \( w^2_i = \left( y^1_i, y^2_i, y^3_i \right) \), and suppose that \( y^2_i < y^1_i < y^3_i < y^1_i \), that is, \( \pi_i^1 = (2, 1, 3, 0) \). To determine the transition probabilities (edge weights as defined in Eq. 6) from \( \pi_i^1 \) to other permutations, we need to find all possible permutations which can be associated to \( \pi_i^2 \) by evaluating the inequalities involving the values in \( w^1_i \) and \( w^2_i \).

We can start by analyzing all possible amplitude relations between \( y^1_i \) and the already established condition \( y^2_i < y^1_i < y^3_i < y^1_i \) to find:

1. \( y^1_i < y^2_i < y^3_i < y^1_i \);
2. \( y^2_i < y^1_i < y^3_i < y^1_i \);
3. \( y^2_i < y^1_i < y^3_i < y^1_i \);
4. \( y^2_i < y^1_i < y^3_i < y^1_i \);
5. \( y^3_i < y^2_i < y^1_i \).

Next, we include \( y^3_i \) and analyze each possible amplitude
relation for all the previous five conditions. This procedure leads to a total of 30 possible amplitude relations between the six elements contained in partitions $w_1$ and $w_2$:

\begin{align}
&i) \ y_1 < y_2 < y_3 < y_1 < y_3 < y_2; \\
&ii) \ y_1 < y_3 < y_2 < y_1 < y_2 < y_3; \\
&\vdots \\
&vii) \ y_2 < y_1 < y_2 < y_1 < y_2 < y_1. \\
\end{align}

By examining the relative positions of $y_1^2, y_1^3, y_2^2, y_2^3$ (which constitute $w_1^2$) in each of the former 30 inequalities, we can ultimately assign an allowed permutation pattern $\Pi$ to $\pi^2$: $i) \Pi = (0, 2, 1, 3); ii) \Pi = (0, 2, 3, 1); \ldots; vii) \Pi = (0, 2, 1, 3); \ldots; xxx) \Pi = (3, 1, 0, 2)$. As we have previously discussed, there are only 12 unique permutations that can horizontally follow $\pi^1$, meaning that some permutations associated with the previous 30 inequalities appear more than once. These different frequencies of occurrence will end-up implying the existence of different edge weights in random ordinal networks.

By repeating the same procedure for $w_1^1 = (y_1^1, y_2^1, y_2^2)$ and $w_2^1 = (y_1^1, y_2^1, y_3^2)$, we find another set of 30 inequalities and their corresponding permutations. Thus, from the analysis of amplitude relations between data values in $w_1^1$ and its neighbors $w_2^1$ and $w_2^2$, we find a total of 60 inequalities corresponding to 60 non-unique permutations. Because amplitude relations in all these inequalities involve random data, all these 60 inequalities are equiprobable [19], and we can count the number of unique permutations (stemming from $\pi^1$) to define their relative frequencies. Finally, we normalize these transition probabilities from $\pi^1$ to all allowed permutations (edge weights Eq. 6) at the node level by dividing the frequency of occurrence of each unique permutation by the total number of possible inequalities involving the elements of $w_1^1$ and its neighbors $w_2^1$ and $w_2^2$ (60 in case $d_x = d_y = 2$). In addition, we also normalize the transitions at the network level by dividing all edge weights by $1/((d_x d_y)!$. This last step is necessary so that the out-strength of permutation $\pi_1^1$ reflects the fact that all different permutations occur with equal probability in random data [17, 24].

We have automatized the former procedure to explicitly consider each possible ordinal pattern in $w_1^1$ (that is, $\pi_1^1 = (0, 1, 2, 3), (0, 1, 3, 2), (0, 2, 1, 3)$, and so on) and thus estimate the edge weights for all allowed transitions in an ordinal network mapped from a large sample of random data. This approach allows us to completely specify a random ordinal network for arbitrary embedding dimensions $(d_x, d_y)$ and estimate all its relevant network metrics (including the global node entropy defined in Eq. 8).

C. Ordinal networks of noisy-periodic ornaments

An interesting application to observe the emergence of random ordinal networks is the randomization process of a geometric ornament [25]. Figure 2(a) illustrates this procedure, where the probability of randomly shuffling pixels values $p$ controls the transition from a periodic image ($p = 0$) to a random one ($p = 1$). We map these sample images (of size $250 \times 250$) into their corresponding ordinal networks for $d_x = d_y = 2$, as shown in Fig. 2(b). The visual aspect of these ordinal networks alludes to the visual features of the noisy-periodic ornaments. For small values of $p$, we observe that the ordinal networks are marked by intense connections between nodes $(1, 0, 3, 2)$ and $(1, 3, 0, 2)$ as well as between $(2, 0, 3, 1)$ and $(2, 3, 0, 1)$, and by auto-loops in these four nodes, which in turn reflect the diagonal stripes in the images. As the values of $p$ increases, the strength of these connections fade out and give rise to a more uniform distribution of edge weights.

To systematically investigate these noisy-periodic ornaments, we generate an ensemble containing 100 ornament samples of size $250 \times 250$ for each randomization probability $p \in \{0.01, 0.05, 0.10, \ldots, 1.00\}$ and transform them into ordinal networks using embedding parameters $d_x = d_y = 2$. From these networks, we estimate the average value of the Gini index of edge weights $G$, the global node entropy $H_{GN}$, and the average weighted shortest path $(l)$, as shown in Figs. 2(c)-(e). In addition, we also estimate the values of these three network properties from the exact form of the random ordinal networks (dashed lines in Figs. 2(c)-(e) presented in Sec. III B).

As expected, we observe that the three measures approach the exact values for random networks with the increase of $p$. It is interesting to notice that the global node entropy $H_{GN}$ surpasses the exact random value around $p = 0.5$ before converging to it. This happens because random ordinal networks are not the most entropic ordinal networks [19], a characteristic that is explained by the fact that edges weights are not all equal in random ordinal networks. We also observe that the average weighted shortest path $(l)$ appears to converge to a value slightly lower than the expected from the exact form of the random ordinal network. This apparent discrepancy occurs because the relatively small size of the images prevents a more accurate estimation of all permutation transitions; however, this difference between the values of $(l)$ for $p = 1$ vanishes as we increase the image size.

D. Ordinal networks of fractional Brownian landscapes

In another application, we investigate ordinal networks mapped from two-dimensional fractional Brownian motion [31]. This class of stochastic processes models natural landscapes and is characterized by the Hurst exponent $h \in (0, 1)$ that controls the surface roughness. Surfaces
FIG. 2. Emergence of random ordinal networks in noisy-periodic ornaments. (a) Visualizations of geometric ornament images for different randomization probabilities \( p \) (shown below images). (b) Ordinal networks with \( d_x = d_y = 2 \) mapped from ornament images displayed in the previous panel. The thicker and darker the edge, the higher is the edge weight (probability associated with permutation transition). Node sizes reflect total in-strength so that the bigger node the more frequently it is found in the symbolic sequence. Dependence of the (c) Gini index of edge weights \( G \), (d) global node entropy \( H_{GN} \), and (e) average weighted shortest paths \( \langle l \rangle \) on the randomization probability \( p \). In the last three panels, the continuous lines show the average values of the network measures and the shaded regions indicate one standard deviation band estimated from an ensemble of 100 ornament samples of size \( 250 \times 250 \) for each value of \( p \). The black dashed lines indicate the exact values of the network measures for random ordinal networks.

generated with small values of \( h \) \((h \to 0)\) are rough while large values of \( h \) \((h \to 1)\) produce smooth landscapes. Cross sections of fractional Brownian landscapes with \( h = 1/2 \) represent usual random walks or Brownian motion. Figure 3(a) shows examples of fractional Brownian landscapes generated by the turning bands method [32] for different values of the Hurst exponent \( h \).

We generate an ensemble containing 100 fractional Brownian landscapes of size \( 256 \times 256 \) for each value of \( h \in \{0.10, 0.15, 0.20, \ldots, 0.90\} \) (with the turning bands method), and map each sample into an ordinal network with embedding dimensions \( d_x = d_y = 2 \). Figure 3(b) presents visualizations of the ordinal networks mapped from the sample images of Fig. 3(a). We observe that changes in surface roughness affect the connectivity patterns of the resulting networks. Rougher surfaces produce ordinal networks with a more even distribution of edge weights which visually resemble random ordinal networks [last panel in Fig. 2(b)]. As fractional Brownian surfaces become smoother, we observe a concentration of weight in a few edges among particular nodes, while the intensity of most links decreases. This concentration of weight reflects the predominant occurrence of only a few permutations in the symbolic arrays related to smoother images.

Beyond the previous qualitative observations, we calculate the average values of the Gini index of edge weights \( G \), the global node entropy \( H_{GN} \), and the average weighted shortest path \( \langle l \rangle \) as a function of the Hurst exponent \( h \) using our ensemble of fractional Brownian landscapes. Figures 3(c)-(e) show these three network measures in comparison with their values estimated from random ordinal networks. These results are in line with our qualitative observations. Specifically, the Gini index...
FIG. 3. Ordinal networks of fractional Brownian landscapes. (a) Examples of fractional Brownian surfaces for a few values of the Hurst exponent (shown below each image). We have normalized all surfaces so that blue shades indicate low height regions and red shades the opposite. (b) Ordinal networks mapped from the fractal surfaces shown in the previous panel. Dependence of the (c) Gini index of edge weights \( G \), (d) global node entropy \( H_{GN} \), and (e) average weighted shortest paths \( l \) on the Hurst exponent \( h \) of two-dimensional fractional Brownian motion. In these last three panels, the continuous lines represent average values (from an ensemble of 100 landscape samples for each \( h \)), and shaded regions delimit an one standard deviation band. The black dashed horizontal lines represent the values of these network metrics estimated from random ordinal networks.

increases with the Hurst exponent, reflecting the concentration of weight in fewer edges. On the other hand, the global node entropy \( H_{GN} \) and the average weighted shortest path \( l \) decrease as the Hurst exponent increases. The monotonic relationships of these three metrics as functions of the Hurst exponent indicate that their values are good predictors of landscape roughness. We have also verified that the behavior of these three network measures obtained for embedding dimensions \( d_x = d_y = 2 \) are very similar to those obtained for \( d_x = 3 \) and \( d_y = 2 \) (or \( d_x = 2 \) and \( d_y = 3 \)). Nevertheless, we need larger fractional Brownian surfaces to properly estimate all transition probabilities.

E. Ordinal networks of Ising surfaces

We have also applied ordinal networks to Ising surfaces [33, 34] to verify whether network measures are capable of identifying phase transitions. These surfaces are obtained by accumulating the spin variables \( \sigma(t) \) of the canonical Ising model in a Monte Carlo simulation [35]. To describe this model, we consider a square lattice whose sites are occupied by spin-1/2 particles \( \sigma \in \{-1, 1\} \) with Hamiltonian given by

\[
\mathcal{H} = -\sum_{\langle i,j \rangle} \sigma_i \sigma_j ,
\]

where the summation is over all pairs of first neighbors. The height \( S_i \) of the corresponding Ising surface at site \( i \) is then defined as

\[
S_i = \sum_t \sigma_i(t) ,
\]

where \( \sigma_i(t) \) is the spin value in step \( t \) of the Monte Carlo simulation.

Figure 4(a) shows examples of Ising surfaces of size 250 × 250 obtained for different reduced temperatures.
$T_r = T/T_c$, where $T_c = 2/\ln(1+\sqrt{2})$ is the critical temperature at which the Ising model undergoes a phase transition. Surfaces generated at reduced temperatures distant from the critical value ($T_r = 1$) do not exhibit long-range structures and are similar to two-dimensional white noise. However, we start to observe more complex patterns as the reduced temperature gets closer, and especially when it is equal, to the critical value.

We generate an ensemble containing 10 Ising surfaces of size $250 \times 250$ for each value of $T_r \in \{0.5, 0.6, \ldots, 2.5\}$. Each surface is obtained after accumulating the spin variables during 30,000 Monte Carlo steps to achieve equilibrium [24, 36]. Next, we map all surfaces into ordinal networks with embedding dimensions $d_x = d_y = 2$. Figure 4(b) shows examples of networks mapped from the images in Fig. 4(a). A visual inspection of these ordinal networks already suggests that their properties change with the reduced temperature.

Similarly to the previous applications, we calculate the Gini index of edge weights $G$, global node entropy $H_{GN}$, and average weighted shortest path ($l$) as a function of the reduced temperature $T_r$. Results presented in Fig. 4(c)-(e) show that all three measures present marked variations before and after the phase transition at $T_r = 1$, assuming extreme values at the critical temperature. The values at the critical temperature are also the furthest from those estimated from random ordinal networks [dashed lines in Fig. 4(c)-(e)]. In addition, we note that variations of these metrics are steeper before than after criticality, indicating that spatial correlations are slowly broken with the rise of temperature. It is also worth observing that these networks become slightly more entropic than random ordinal networks at high temperatures.

**F. Ordinal networks of texture images**

As a final application, we use our generalized algorithm to map real-world images to ordinal networks. To do so, we consider 112 pictures of natural textures comprising the Normalized Brodatz Texture database [37], a set of...
FIG. 5. Ordinal networks mapped from Brodatz textures. (a) Locations of all Brodatz textures at the plane of global node entropy $H_{GN}$ versus average weighted shortest path $(l)$. (b) Six different images corresponding to highlighted textures in the previous panel (blue markers). (c) Difference between the global network entropy estimated from horizontal and vertical ordinal networks ($H_{GN}^{\text{Horizontal}} - H_{GN}^{\text{Vertical}}$). (d) Six pictures corresponding to highlighted textures in the previous panel (red markers).

8-bit (256 gray levels) images of size $640 \times 640$ covering a large variety of natural textures that is often used as validation set for texture analysis [37]. We map each image into an ordinal network with embedding dimensions $d_x = d_y = 2$. We then estimate the Gini index of edge weights $G$, global node entropy $H_{GN}$, and average weighted shortest path $(l)$ of each network.

These network measures are cross-correlated with Pearson correlation coefficients ranging from 0.92 for $(l)$ versus $H_{GN}$ to $-0.96$ for $G$ versus $(l)$. Figure 5(a) shows a scatter plot of $(l)$ versus $H_{GN}$ (the less significant association) for all textures, where the dispersion pattern suggests that both measures are non-linearly related. Figure 5(b) depicts six different images that are also highlighted in Fig. 5(a). We observe that the two textures with extreme values of $(l)$ and $H_{GN}$ (D49 and D71) are quite different: while regular horizontal stripes mark texture D49, texture D71 shows much more complex structures.

We also note the existence of textures with similar values of $H_{GN}$ and distinct values of $(l)$ as well as images with similar values of $(l)$ and different values of $H_{GN}$. These results suggest that both measures may quantify different aspects of images. For instance, textures D86 and D81 [highlighted in Fig. 5(a)] have almost the same values of global node entropy; however, the value of $(l)$ is considerably larger for D86 than D81. By inspecting these two textures in Fig. 5(b), we note that D86 appears to be rougher than D81. Similarly, textures D18 and D78 have comparable values of $(l)$ but quite different values of $H_{GN}$. The visual inspection of these textures suggests that D18 is more regular and structured than D78. While it is challenging to generalize these interpretations to other images, we believe the values of $H_{GN}$ quantify patterns at a more local level while $(l)$ and $G$ are more global measures. This idea somehow agrees with the definition of these measures in the sense that $H_{GN}$ is based on relations involving first-neighbors, while $(l)$ and $G$ involve the entire ordinal network.

We further investigate the possibility of exploring visual symmetries in the Brodatz data set. To do so, we have made a small modification in our original algorithm to create two ordinal networks from a single image. One of these networks considers only horizontal transitions among permutations (horizontal ordinal network), and the other uses solely vertical transitions among permutations (vertical ordinal network). Thus, we map each Brodatz texture into a horizontal and vertical ordinal network with embedding dimensions $d_x = d_y = 2$. Then, we estimate the global node entropy from the horizontal network ($H_{GN}^{\text{Horizontal}}$) and the vertical network ($H_{GN}^{\text{Vertical}}$). Figure 5(c) shows the difference between these quantities ($H_{GN}^{\text{Horizontal}} - H_{GN}^{\text{Vertical}}$) for each texture. We observe a few textures with extreme values for this difference and highlight six of them, which are also depicted in Fig. 5(d). Most of these images are characterized by stripes or line segments predominantly oriented in vertical or horizontal directions, indicating that vertical and horizontal ordinal
networks are capable of detecting this simple symmetry feature.

These results indicate potential applications of our approach in classification tasks as a way of extracting texture features via network metrics [23], in a similar manner to which permutation based metrics have already been successfully used as predictors in image classification and regression problems [26, 27].

IV. CONCLUSION

We have proposed a generalization of the ordinal network algorithm for mapping images (two-dimensional data) into networks. After describing the method, we have studied basic connectivity patterns of these networks which in turn allowed us to find the exact form of ordinal networks mapped from random data. We have observed the emergence of these random ordinal networks in a controlled setting by randomizing a geometric periodic ornament. We have also investigated changes in surface roughness of two-dimensional fractional Brownian motion and found monotonic relations between network measures and the Hurst parameter of these fractal surfaces. This result extends similar findings previously obtained from time series data [19]. In the context of physical models, we have studied phase transition in an Ising-like model where variations in network metrics have been proven useful for accurately identifying the critical temperature. Finally, we have mapped natural texture images into ordinal networks and briefly considered the possibility of using networks to potentially extract image features.

Our work thus contributes to recent developments in network science focused on mapping two-dimensional data into networks and characterizing them through network representations [21, 22]. As these network approaches are quite novel when compared to other attempts stemming from complexity science [24, 25, 36, 38–40], we believe there are several opportunities for applying these new tools to different contexts involving image analysis.

[1] A. Vespignani, Nature 558, 528 (2018).
[2] S. Dorogovtsev, Lectures on Complex Networks, Oxford Master Series in Physics (Oxford University Press, 2010).
[3] R. Albert and A.-L. Barabási, Reviews of Modern Physics 74, 47 (2002).
[4] M. Zanin, D. Papo, P. Sousa, E. Menasalvas, A. Nicchi, E. Kubik, and S. Boccaletti, Physics Reports 635, 1 (2016).
[5] Y. Zou, R. V. Donner, N. Marwan, J. F. Donges, and J. Kurths, Physics Reports 787, 1 (2019).
[6] L. Lacasa, B. Luque, F. Ballesteros, J. Luque, and J. C. Núñez, Proceedings of the National Academy of Sciences 105, 4972 (2008).
[7] N. Marwan, J. F. Donges, Y. Zou, R. V. Donner, and J. Kurths, Physics Letters A 373, 4246 (2009).
[8] M. Small, in 2013 IEEE International Symposium on Circuits and Systems (ISCAS2013) (2013) pp. 2509–2512.
[9] B. Luque, L. Lacasa, F. Ballesteros, and J. Luque, Physical Review E 80, 046103 (2009).
[10] L. Lacasa, A. Núñez, É. Roldán, J. M. R. Parrondo, and B. Luque, The European Physical Journal B 85, 217 (2012).
[11] I. Bezsudnov and A. Snarskii, Physica A 414, 53 (2014).
[12] F. M. Bianchi, L. Livi, C. Alippi, and R. Jenssen, Scientific Reports 7, 44037 (2017).
[13] J.-P. Eckmann, S. O. Kamphorst, and D. Ruelle, Europhysics Letters (EPL) 4, 973 (1987).
[14] H. Kantz and T. Schreiber, Nonlinear Time Series Analysis, Cambridge Nonlinear Science Series (Cambridge University Press, 2004).
[15] E. Bradly and H. Kantz, Chaos 25, 097610 (2015).
[16] C. Bandt and F. Shiha, Journal of Time Series Analysis 28, 646 (2007).
[17] C. Bandt and B. Pompe, Physical Review Letters 88, 174102 (2002).
[18] M. McCullough, M. Small, T. Stemler, and H. H.-C. Ih, Chaos 25, 053101 (2015).
[19] A. A. B. Pessa and H. V. Ribeiro, Physical Review E 100, 042304 (2019).
[20] F. Olivares, M. Zanin, L. Zunino, and D. Pérez, Chaos 30, 063101 (2020).
[21] Q. Xiao, X. Pan, X.-L. Li, M. Stephen, H.-J. Yang, Y. Jiang, J.-Y. Wang, and Q.-J. Zhang, Chinese Physics B 23, 078904 (2014).
[22] L. Lacasa and J. Iacovacci, Physical Review E 96, 012318 (2017).
[23] J. Iacovacci and L. Lacasa, IEEE Transactions on Pattern Analysis and Machine Intelligence 42, 974 (2019).
[24] H. V. Ribeiro, L. Zunino, E. K. Lenzi, P. A. Santoro, and R. S. Mendes, PLOS ONE 7, 1 (2012).
[25] L. Zunino and H. V. Ribeiro, Chaos, Solitons & Fractals 91, 679 (2016).
[26] H. Y. D. Sigaki, R. F. de Souza, R. T. de Souza, R. S. Zola, and H. V. Ribeiro, Physical Review E 99, 013311 (2019).
[27] H. Y. D. Sigaki, M. Perc, and H. V. Ribeiro, Proceedings of the National Academy of Sciences 115, E8585 (2018).
[28] Y. Cao, W.-w. Tung, J. B. Gao, V. A. Protopopescu, and L. M. Hively, Physical Review E 70, 046217 (2004).
[29] O. A. Rosso, H. A. Larrondo, M. T. Martin, A. Plastino, and M. A. Fuentes, Physical Review Letters 99, 154102 (2007).
[30] M. McCullough, M. Small, H. H. C. Ih, and T. Stemler, Philosophical Transactions of the Royal Society A 375, 20160292 (2017).
[31] B. B. Mandelbrot, The Fractal Geometry of Nature (Freeman, San Francisco, 1982).
[32] Z.-M. Yin, Journal of Computational Physics 127, 66 (1996).
[33] A. F. Brito, J. A. Redinz, and J. A. Plascak, Physical Review E 75, 046106 (2007).
[34] A. F. Brito, J. A. Redinz, and J. A. Plascak, Physical Review E 81, 031130 (2010).
[35] D. Landau and K. Binder, A Guide to Monte Carlo Simulations in Statistical Physics (Cambridge University Press, 2015).
[36] A. Brazhe, Physical Review E 97, 061301 (2018).
[37] A. Safia and D.-C. He, ISRN Machine Vision 2013, 14 (2013).
[38] Z. Cai, E. Shen, F. Gu, Z. Xu, J. Ruan, and Y. Cao, International Journal of Bifurcation and Chaos 16, 3235 (2006).
[39] Y. Andrienko, N. Brilliantov, and J. Kurths, The European Physical Journal B 15, 539 (2000).
[40] D. P. Feldman and J. P. Crutchfield, Physical Review E 67, 051104 (2003).