The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

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Abstract. Motivated by recent studies by Dorey, Pocklington and Tateo for unitary minimal models perturbed by $\phi_{1,2}$, we examine the thermodynamics of one-dimensional quantum systems, whose counterparts in the two-dimensional classical model are the dilute $A_L$ models in regime 2. The functional relations for arbitrary values of $L$ are established. Guided by numerical evidence, we obtain a set of coupled integral equations from the established relations, which yields the evaluation of the free energy at arbitrary temperature. In the scaling limit, the integral equations coincide with the thermodynamic Bethe ansatz equations (TBA) proposed by Dorey, Pocklington and Tateo, thereby supporting their results. The new fermionic representations of the Virasoro characters are remarked upon briefly.

Keywords: integrable quantum field theory, integrable spin chains (vertex models), solvable lattice models, thermodynamic Bethe ansatz

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Contents

1. Introduction 3

2. The dilute $A_L$ models and the quantum transfer matrix 5

3. The fusion transfer matrices and quantum Jacobi–Trudi formula 9
   3.1. The symmetric fusion structure .................................. 10
   3.2. The quantum Jacobi–Trudi formula ............................. 13

4. Breather $T$-system ......................................................... 15

5. The $sl_3$ type functional relations .................................... 17

6. The magnon-like $t$-system and the Dynkin-like diagrams 19
   6.1. The magnon-like $t$-system for $A_{4k-1}$ .......................... 20
   6.2. The magnon-like $t$-system for $A_{4k}$ ............................ 21
   6.3. The magnon-like $t$-system for $A_{4k+1}$ ........................ 22
   6.4. The magnon-like $t$-system for $A_{4k+2}$ ........................ 23
   6.5. The solutions to the magnon-like $t$-system ..................... 23

7. Kink transfer matrix ......................................................... 23

8. The reality property of magnon-like $t$-functions .............. 25

9. The $Y$-system and the $T$-system ........................................ 25

10. Proof of the magnon-like $t$-system for $L = 4k+2$ ............... 27

11. Thermodynamic Bethe ansatz in the scaling limit ............... 29
   11.1. Analytic properties .................................................. 29
   11.2. The coupled integral equations .................................... 30
   11.3. Thermodynamic Bethe ansatz equations ......................... 32

12. Summary and discussions .................................................. 34

Acknowledgments ................................................................. 35

Appendix A. $Y$-system ......................................................... 35
   Appendix A.1. Case $(L = 4k - 1, k \geq 2)$ ......................... 36
   Appendix A.2. Case $(L = 4k, k \geq 2)$ ............................... 36
   Appendix A.3. Case $(L = 4k + 1, k \geq 2)$ ......................... 37
   Appendix A.4. Case $(L = 4k + 2, k \geq 2)$ ......................... 38

Appendix B. Explicit solutions to magnon-like $t$-system .......... 38
   Appendix B.1. Solutions for $A_{L=4k-1}$ ............................ 38
   Appendix B.2. Solutions for $A_{L=4k}$ ................................. 40
   Appendix B.3. Solutions for $A_{L=4k+1}$ ............................ 41
   Appendix B.4. Solutions for $A_{L=4k+2}$ ............................ 43

Appendix C. Functional relations for the kink transfer matrix ..... 43

Appendix D. Proofs of lemmas 5 and 6 ..................................... 45

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Appendix E. Proof of the magnon-like t system for $L = 4k + 1$

Appendix E.1. The proof for $t_1^{(3)}(x \pm \frac{4}{2L+4}i) = t_1^{(2)}(x) + t_2^{(1)}(x)$ .... 48

Appendix E.2. The proof for $t_2^{(3)}(x \pm \frac{1}{2L+4}i) = t_1^{(2)}(x)t_1^{(7)}(x) + t_2^{(2)}(x)$ .... 48

Appendix E.3. The proof for $t_1^{(3)}(x \pm \frac{4}{2L+4}i) = t_1^{(1)}(x)t_1^{(5)}(x) + t_2^{(3)}(x)$ .... 48

Appendix E.4. The proof for $t_2^{(3)}(x \pm \frac{1}{2L+4}i) = t_1^{(1)}(x \pm \frac{1}{2L+4})t_2^{(5)}(x) + t_1^{(3)}(x)t_3^{(3)}(x)$ .... 49

Appendix E.5. The proof for $t_1^{(6)}(x \pm \frac{1}{2L+4}i) = t_1^{(5)}(x) + t_2^{(4)}(x)$ .... 49

Appendix E.6. The proof for $t_1^{(4)}(x \pm \frac{1}{2L+4}i) = t_1^{(2)}(x)t_1^{(5)}(x) + t_2^{(4)}(x)$ .... 50

Appendix E.7. The proof for $t_2^{(4)}(x \pm \frac{1}{2L+4}i) = t_1^{(2)}(x \pm \frac{1}{2L+4})t_2^{(5)}(x) + t_1^{(4)}(x)t_3^{(4)}(x)$ .... 51

Appendix E.8. The proof for $t_1^{(5)}(x \pm \frac{1}{2L+4}i) = t_1^{(4)}t_1^{(6)} + t_2^{(5)}(x)$ .... 53

Appendix F. A result for an exceptional case

Appendix G: Numerical results for the dilute $A_L$ model in regime 2

Appendix H. The kernel matrices for $L = 5, 7$

References

1. Introduction

The Andrews–Baxter–Forrester model [1] has been a prototype in the studies of integrable lattice models as well as the integrable $\Phi_{1,2}$ perturbation of the minimal unitary CFT.

The dilute $A_L$ model proposed in [2,3] deserves the same attention, as it represents a lattice analogue of the $\phi_{1,2}$ and $\phi_{2,1}$ perturbations of the minimal unitary CFT.

Especially, there are abundant research results on the dilute $A_L$ model in regime 2, which corresponds to the model of high interest, $M_{3,4} + \Phi_{1,2}$ theory. The latter belongs to the same universality class as the Ising model in a magnetic field at $T_c$, and is characterized by the famous $E_8$ structure [4–12]. There is much evidence for this coincidence, the central charge [2,3], surface exponents [13], universal amplitude ratio [14–16], scaling dimensions [17] and excitation spectra [18]. The most impressive demonstration may be presented in [19], which confirms the hidden $E_8$ structure behind the thermodynamic Bethe ansatz equation (TBA, for short) of the dilute $A_3$ model. See also [20,21].

Similarly, $M_{4,5} + \Phi_{1,2}$ ($M_{6,7} + \Phi_{1,2}$) theory has $E_7$ ($E_6$) structure hidden behind [4]. In spite of the lack of a relevant string hypothesis, the technique of the quantum transfer matrix (QTM) and the fusion relations make it possible to manifest the desired Lie algebraic structure in the TBA of the dilute $A_{3,4,6}$ in regime 2 [22,23].

The aim of the present report is to extend the result further; we will establish a closed set of fusion relations of the dilute $A_L$ model for arbitrary $L$. From these relations, we derive the nonlinear integral equations (TBA) which yield the evaluation of the free energy of the corresponding 1D quantum model at arbitrary temperatures. This is obviously motivated by the recent progress in the study of the models as perturbed conformal field theories [24,25]. The perturbation theory by an operator $\Phi_{1,2}$ or $\Phi_{2,1}$ has encountered serious technical difficulties. The systematic studies on the bootstrap procedure on the
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

S-matrix have been initiated in [26] and [27]. See also [28,29]. The latter approach is based on the scaling $q$-state Potts field theory and will be relevant in our view. It is only recently that the closed bootstrap procedure, associated with the Potts field theory, has been obtained for a set of $S$-matrices for a wide range of parameters by Dorey et al [24]. The check of the results against a finite size system however again poses a challenging problem because of the non-diagonal nature of the scattering theory. In [25], a set of involved TBA equations is conjectured by ingenious insights from special cases which have similarity to the TBA for the sine–Gordon model. They prepare pieces of TBA equations, one of which comes from $Z_n$ models [30], then glue them together by respecting symmetries and by physical requirements. The resultant equations pass several non-trivial checks, but lack firm grounds. It is thus interesting to see if it is possible to recover the $Y$-system [31]–[33], the algebraic version of the TBA, from the fusion relations of the dilute $A_L$ model.

The functional relations among transfer matrices (the $T$-system) are determined by the fusion relations [34,36]. The decomposition of the product of two modules at the singular points of the Boltzmann weight rules the latter. Thus the problem seems to be straightforward at first sight. This is, unfortunately, not true. There are often infinitely many functional relations for a given model. One thus has to choose a subset which offers the closed $T$-system with finitely many entries. Moreover, among the sets, only those of which both sides contain fusion transfer matrices with good analytic properties (to be specified) will be relevant for the TBA. Therefore, the problem of finding the relevant functional relations so as to reach the TBA is far from trivial.

It is well known that the dilute $A_L$ model has $U_q(A_{2}^{(2)})$ symmetry at the critical point. The off-critical extension, deformed Virasoro algebra $\mathcal{V}_{x,r}(A_{2}^{(2)})$, plays a role in the vertex operator approach [37]. This elegant result, however, is not of direct relevance. The crucial fact for the present study is, rather, that the fusion structure is valid even away from criticality [38]; it possesses fusion structures similar to $U_q(A_{1}^{(1)})$ and $U_q(A_{2}^{(1)})$. We will relate the former to breathers while the latter to the pseudo-particles constituting ladders in their TBA diagrams. In addition, we will introduce a set of objects which we call ‘magnon-like’, characterized by similar diagrams to those associated with simply laced Lie algebra ($A, D$ and $E$). It turns out that they constitute closed functional relations with good analyticity. They are three almost independent relations but for one common object $T_{B_3}$, which plays the role of glue. Thus the construction employed here makes the meaning of gluing in [25] very explicit.

We will introduce auxiliary functions, fusion transfer matrices\(^1\), labelled by skew Young diagrams [39]–[42]. In view of the latter diagrams, proofs of the functional relations seem to be similar to those of the decomposition rules in the symmetric group. There are, however, extra complications due to the fact that our coupling constant in a sense lies at a ‘root of unity’. Sometimes two (or even more) diagrams, with completely different appearances, correspond to an identical fusion transfer matrix. One must choose relevant expressions (diagrams) case by case in order to obtain suitable decompositions. This makes the proof of the functional relations simple but lengthy, unfortunately. There are

\(^1\) We expect them to be the eigenvalues of transfer matrices of which fusion types are described by the diagrams. This is however yet to be proved. As our argument does not rely on whether this coincidence is true or not, we simply call the auxiliary functions (eigenvalues of) fusion transfer matrices.
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Figure 1. An incidence diagram for the dilute $A_3$ model. The local states corresponding to connected nodes can be located to nearest neighbour sites on a square lattice.

four families of functional relations depending on the values of $L$ and they need separate treatment at the present stage.

The validity of the obtained algebraic equations does not necessary legitimate the integral equations (TBA). We need another step to clarify the analytic structure of $Y$-functions. This remains an important issue to be understood. In this report, we employ the direct numerical calculation by various ranges of parameters and adopt a simple conjecture on the analyticity. With the established functional relations and the assumption on the analyticity, it is then straightforward to transform the $T$-system into the $Y$-system and then the $Y$-system into coupled integral equations [35]. In the scaling limit, the TBA proposed in [25] is recovered from them. A simple ansatz, suggested by numerics, on the zeros of the fusion transfer matrices remarkably leads to the correct mass ratio of the particles determined by the $S$-matrix poles.

This paper is organized as follows. In the next section, we give a brief review on the dilute $A_L$ models and the QTM method. Some obvious consistency of the lattice models and the CFT perturbed by $\Phi_{1,2}$ will be remarked. Section 3 is devoted to the discussion on fusion transfer matrices which are parametrized by skew Young diagrams. We introduce three important kinds of objects in sections 4–6. The first object introduced in section 4 seems to correspond to the breathers in [24] while the second one in section 5 is responsible for the ladder structure in the $Y$-system diagram in [25]. Section 6 is divided into four subsections; each is devoted to the discussion on ‘magnon-like’ objects. In the case of the dilute $A_L$ model, $L$ even, a fundamental role seems to be played by a ‘kink’ transfer matrix, which is introduced in section 7. We will give a remark on the ‘reality property’ of the magnon-like objects in section 8. Once fusion relations, the $T$-system, are established, it is readily transformed into the $Y$-system. The explicit procedure will be discussed in section 9. The proofs of the $T$-system (or the ‘magnon-like’ $t$-system) are simple but lengthy. Those for $L$ even (odd) are somewhat similar. For brevity, we thus present the proof for $L = 4k + 2$ in section 10 and $L = 4k + 1$ in appendix F. In section 11, it will be shown that the conjectured TBA is naturally restored in a scaling limit. We conclude the paper with a brief summary and discussion in section 12. A new fermionic formula for the Virasoro character found for $M_{5,6}$ is briefly commented.

2. The dilute $A_L$ models and the quantum transfer matrix

The dilute $A_L$ model is proposed in [2] as an elliptic extension of the Izergin–Korepin model [43]. The model is of the restricted SOS type with local variables $\in \{1, 2, \ldots, L\}$. 

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

The variables $\{a, b\}$ on neighbouring sites should satisfy the adjacency condition, $|a - b| \leq 1$, which is often described by a graph in figure 1. The following RSOS weights are found in [2] which satisfy the Yang–Baxter relation. They are parametrized by the spectral parameter $u$, the elliptic nome $q = e^{-\pi \tau}$. 

\[
a \begin{bmatrix} a \\ u \\ a \end{bmatrix} = \begin{bmatrix} \theta_1(6 - u)\theta_1(3 + u) \\ - \theta_1(u)\theta_1(3 - u) \end{bmatrix} \begin{bmatrix} S_{a+1} \theta_4(2a - 5) \\ S_a \theta_4(2a + 1) \end{bmatrix} + \frac{S_{a-1} \theta_4(2a + 5)}{S_a \theta_4(2a - 1)},
\]

\[
a \pm 1 \begin{bmatrix} a \\ u \\ a \end{bmatrix} = \pm \begin{bmatrix} \theta_1(3 - u)\theta_4(\pm 2a + 1 - u) \\ \theta_1(3 + u)\theta_4(\pm 2a + 1) \end{bmatrix},
\]

\[
a \pm 1 \begin{bmatrix} a \\ u \\ a \end{bmatrix} = \begin{bmatrix} \theta_1(3 - u)\theta_4(\pm 2a + 2 + u) \\ \theta_1(3 + u)\theta_4(\pm 2a + 2) \end{bmatrix},
\]

\[
a \pm 1 \begin{bmatrix} a \\ u \\ a \end{bmatrix} = \begin{bmatrix} \theta_1(3 - u)\theta_4(\pm 4a + 2 + u) \\ \theta_1(3 + u)\theta_4(\pm 4a + 2) \end{bmatrix} + \frac{S_{a+1} \theta_1(u)\theta_1(\pm 4a - 1 + u)}{S_a \theta_1(\pm 4a + 2)},
\]

\[
a \pm 1 \begin{bmatrix} a \\ u \\ a \end{bmatrix} = \begin{bmatrix} \theta_1(3 - u)\theta_4(\pm 4a - 4 + u) \\ \theta_1(3 + u)\theta_4(\pm 4a - 4) \end{bmatrix} + \frac{S_{a+1} \theta_1(u)\theta_1(\pm 4a - 5)}{S_a \theta_1(\pm 4a + 2)} \theta_1(1)\theta_1(\pm 4a + 4 + u),
\]

Here $\theta_{1,4}(x) = \vartheta_{1,4}(\lambda x, \tau)$,

\[
\vartheta_1(x, \tau) = 2q^{1/4} \sin x \prod_{n=1}^\infty (1 - 2q^{2n} \cos 2x + q^{4n})(1 - q^{2n}),
\]

\[
\vartheta_4(x, \tau) = \prod_{n=1}^\infty (1 - 2q^{2n-1} \cos 2x + q^{4n-2})(1 - q^{2n}),
\]

\[
S_a = (-1)^a \frac{\theta_1(4a)}{\theta_1(2a)}
\]

Note an important property, the standard initial condition, that only those weights with identical variables on the NE and the SW diagonal remain non-zero when $u \to 0$.

The crossing parameter $\lambda$ needs to be a function of $L$ for the restriction. The model exhibits four different physical regimes depending on parameters [2, 3]:

- **regime 1**, $0 < u < 3$, $\lambda = \frac{\pi L}{4(L+1)}$, $L \geq 2$;
- **regime 2**, $0 < u < 3$, $\lambda = \frac{\pi(L+2)}{4(L+1)}$, $L \geq 3$;
- **regime 3**, $3 - \frac{\pi}{\lambda} < u < 0$, $\lambda = \frac{\pi(L+2)}{4(L+1)}$, $L \geq 3$;

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

- **regime 4**, $3 - \frac{\pi}{\lambda} < u < 0$, $\lambda = \frac{\pi L}{4(L+1)}$, $L \geq 2$.

In the low temperature limit, the ground states correspond to ferromagnetic ordered states in regimes 1 and 2 in which we are interested.

As usual, we define the row to row transfer matrix $T_{RTR}(u)$ by

$$(T_{RTR}(u))_{\{b\}}^{\{a\}} = \prod_{j=1}^{M} b_j a_j u^{b_j+1} a_{j+1}^{b_{j+1}}.$$  

Denoting its largest eigenvalue by the same symbol, $T_{RTR}(u)$, the free energy per site of the 2D classical model in the thermodynamic limit reads $\beta f = -\lim_{M \to \infty} \frac{1}{M} \log T_{RTR}(u)$. Several other quantities other than the free energy have been evaluated for the models. For example, we mention the explicit one-point functions elaborated in [17].

We are not dealing with the dilute $A_L$ model directly but its one-dimensional quantum counterpart.

The Hamiltonian for the latter is defined by

$$\mathcal{H} = -2 \frac{\partial}{\partial u} \ln T_{RTR}(u) \bigg|_{u=0}$$

as in [19]. The explicit form of the Hamiltonian is not needed in the following argument.

We are interested in the free energy of the 1D model at finite temperatures. The success in the previous studies on the thermodynamics of 1D quantum systems depends on the efficient conjectures on the dominant solutions of the Bethe ansatz equation in the thermodynamic limit, the string hypothesis. It seems too difficult to perform the same procedure for arbitrary $L$. Instead, we apply different machinery, the method of the quantum transfer matrix (QTM) [44]–[48].

The quantity of interest is the partition function of the 1D model,

$$Z_{1D} = \text{Tr} e^{-\beta \mathcal{H}}.$$  

The first step is to rewrite this as the partition function of the 2D model, defined on an $M \times N$ square lattice by a simple trick,

$$Z_{1D} = \lim_{N \to \infty} \text{Tr} \left( \exp \left( -\frac{\beta}{N} \mathcal{H} \right) \right)^{N} = \lim_{N \to \infty} \text{Tr} (T^\vee)^{N} = \lim_{N \to \infty} \text{Tr} Z_{2D}. \quad (3)$$

The fictitious dimension, $N$, is referred to as the Trotter number. $T^\vee$ is a kind of transfer matrix, composed of matrix elements of $e^{-\frac{\beta}{N} \mathcal{H}}$.

In order to relate $T^\vee$ to the row to row transfer matrix above, we introduce a rotated transfer matrix, $\tilde{T}_{RTR}(u)$,

$$(\tilde{T}_{RTR}(u))_{\{b\}}^{\{a\}} = \prod_{j=1}^{M} b_j a_j \Xi^{b_j+1} a_{j+1}^{b_{j+1}}.$$  

Thanks to (2) and the standard initial condition, we immediately see

$$T_{RTR}(u) \sim P_R \exp \left( -\frac{u}{2} \mathcal{H} + O(u^2) \right), \quad \tilde{T}_{RTR}(u) \sim P_L \exp \left( -\frac{u}{2} \mathcal{H} + O(u^2) \right),$$

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where \( P_R \) (\( P_L \)) is the right (left) shift operator, commuting with the Hamiltonian and \( P_L P_R = P_R P_L = \text{the identity operator} \). Then the desired expression for \( T^\vee \) is obtained,

\[
T^\vee = T_{RTR} \left( \frac{\beta}{N} \right) T_{RTR} \left( \frac{\beta}{N} \right)
\]

where the limit \( N \to \infty \) is assumed on both sides. The eigenvalues of \( T^\vee \) are highly degenerated in \( N \to \infty \), which makes the summation in the partition function difficult to perform. The intriguing observation \([44]\) is that if one defines yet another transfer matrix \( T \), the quantum transfer matrix (QTM), propagating in the horizontal direction, then there is a gap in between the largest and the second largest eigenvalues which remains non-vanishing in the limit \( N \to \infty \). Now that \( Z_{1D} \sim \text{Tr} T^M \) and \( M \to \infty \), the only largest eigenvalue of \( T \) is necessary in the evaluation of \((3)\).

In the most sophisticated formulation \([49]\), the QTM is defined with one further variable \( x \),

\[
(T_{\text{QTM}}(u, x))_{\{a\}}^{\{b\}} = \prod_{j=1}^{N/2} \begin{vmatrix}
 b_{2j-1} & a_{2j-1} & u + ix \\
 b_{2j} & a_{2j} & u \\
 b_{2j+1} & a_{2j+1} & u + ix
\end{vmatrix}.
\]

By taking \( x = 0 \) one recovers \( T \).

The commutative property of QTMs with fixed \( u \) is crucial in the QTM formulation,

\[
[T_{\text{QTM}}(u, x), T_{\text{QTM}}(u, x')] = 0.
\]

It represents the existence of infinitely many conserved quantities commuting with the QTM, which may be a key in the exact enumeration of the free energy at arbitrary temperatures. Since we will adopt an identical value for \( u = \frac{\beta}{N} \) for commuting the QTM, it will be dropped afterwards.

Let us explicitly write again that the free energy per site is represented only by the largest eigenvalue of \( T_{\text{QTM}}(x = 0) \) with \( u = \frac{\beta}{N} \),

\[
\beta f = - \lim_{M \to \infty} \frac{1}{M} \ln \text{Tr} \exp(-\beta H) = - \lim_{N \to \infty} \ln (\text{the largest eigenvalue of } T_{\text{QTM}}(x = 0)).
\]

Our aim is thus to evaluate the largest eigenvalue of the QTM. This is again not such a simple problem as it seems at first sight. The coupling constant \( u \) depends on the fictitious system size, the Trotter number, which makes it difficult to take the limit \( N \to \infty \).

Our strategy is to introduce auxiliary transfer matrices (fusion transfer matrices) commuting with the QTM and to explore functional relations among them \([22, 23], [50]–[53]\). We find the latter relations, if a suitable subset is chosen, are efficient enough to yield the evaluation of the largest eigenvalue of \( T_{\text{QTM}} \) for arbitrary \( N \): the functional relations can be transformed into coupled integral equations whose solution yields the desired quantity. The resultant equations allow us to take the limit \( N \to \infty \) without any difficulty.

In the next section, we will introduce the fusion transfer matrices. Before doing so, we mention simple evidence that the dilute \( A_L \) model (or its 1D analogue) is a lattice analogue of the minimal unitary model perturbed by the \( \Phi_{1,2} \) or \( \Phi_{2,1} \) operator. In TBA formulation, the dressed energy function \( \epsilon \) and \( Y = \exp(\epsilon) \) are main objects. It is pointed out in \([54]\)
that, within the framework of TBA, it seems too difficult to show analytically that the small $r$ expansion coincides with the result from the conformal perturbation theory. The hidden periodicity of the $Y$-system, found by a brute-force substitution, partially resolves the problem [31].

We will show, within our formalism, that $Y$ is expressible by ratios of fusion transfer matrices. Therefore, in our view, the periodicity of $Y$ (and $\epsilon$) is a consequence of the periodicity of the original Boltzmann weights. In view of the rapidity variable $\theta \equiv \frac{4\pi}{\lambda}$, the periods read

$$P_\theta = \begin{cases} 
\frac{4(L+1)\pi}{3L}, & \text{for regime 1} \\
\frac{4(L+1)\pi}{3(L+2)}, & \text{for regime 2}.
\end{cases} \tag{4}$$

From this one can read off the scaling dimension of the perturbed field by $\Delta = 1 - \frac{\pi}{P_\theta}$, explicitly,

$$\Delta = \begin{cases} 
\frac{p+3}{4p}, & \text{for regime 1} \\
\frac{p-2}{4(p+1)}, & \text{for regime 2}.
\end{cases} \tag{5}$$

Note that we adopted a different identification,

$$p = \begin{cases} 
L+1, & \text{for regime 1} \\
L, & \text{for regime 2}.
\end{cases} \tag{6}$$

They coincide with the scaling dimensions $\Delta_{2,1}$ and $\Delta_{1,2}$ in the unitary model $M_{p,p+1}$. This thus supports the hypothesis that the dilute $A_L$ models in regime 1 (regime 2) are lattice regularizations of perturbed minimal models $M_{p,p+1}$ under the identification (6).

From now on, we restrict our argument to regime 2. As is announced, the fusion transfer matrices will be examined in the next section.

3. The fusion transfer matrices and quantum Jacobi–Trudi formula

Any model integrable in the sense of the Yang–Baxter relation offers a systematic way of generalization, the fusion procedure. It utilizes the singular values of spectral parameter at which the Boltzmann weights reduce to projectors to a subspace. In [38], a detailed description is given for the explicit procedures to obtain ‘generalized’ Boltzmann weights.

The explicit construction of these weights for various fusion types is not of issue in this report. Rather, the main objects of our interest are functional relations which involves QTM. In the following we introduce several auxiliary functions which generalizes QTM. We expect that they correspond to eigenvalues of transfer matrices made by the fusion procedure. However, we do not attempt to prove this, as it is not necessary for the evaluation of free energy.
The most important building block in our argument is the following expression for the (any) eigenvalues of QTM, which we will denote by $T_1(x)$,

$$T_1(x) = w \phi \left( x + \frac{3}{2} \right) \phi \left( x + \frac{1}{2} \right) Q(x - \frac{3}{2}i) Q(x + \frac{3}{2}i) + w^{-1} \phi \left( x - \frac{3}{2}i \right) \phi \left( x - \frac{1}{2}i \right) Q(x + \frac{3}{2}i) Q(x + \frac{1}{2}i)$$

$$+ w^{-1} \phi \left( x - \frac{3}{2}i \right) \phi \left( x - \frac{1}{2}i \right) Q(x + \frac{3}{2}i) Q(x + \frac{1}{2}i),$$

(7)

$$Q(x) := \prod_{j=1}^{N} h[x - x_j],$$

$$\phi(x) := \left( \frac{h[x + (\frac{3}{2} - u)i]h[x - (\frac{3}{2} - u)i]}{h[2i]h[3i]} \right)^{N/2}, \quad h[x] := \theta_1(ix),$$

where $w = \exp(i\frac{\pi \ell}{L+1})$ and $\ell = 1$ for the largest eigenvalue sector. It is often referred to as the dressed vacuum form.

The parameters, $\{x_j\}$, are solutions to the ‘Bethe ansatz equation’ (BAE) [19],

$$w \frac{\phi(x_j + i)}{\phi(x_j - i)} = - \frac{Q(x_j - i)Q(x_j + 2i)}{Q(x_j + i)Q(x_j - 2i)}, \quad j = 1, \ldots, N.$$  

(8)

We introduce

$$P = \begin{cases} 
\frac{4(L+1)}{L}, & \text{for regime 1} \\
\frac{4(L+1)}{L+2}, & \text{for regime 2}. 
\end{cases}$$

(9)

Then the periodicity $T_1(x) = T_1(x + Pi)$ follows obviously from the above expression.

We represent the three terms in the eigenvalue of the quantum transfer matrix by three boxes with letters 1, 2 and 3 suffixed by the spectral parameter.

$$T_1(u, x) = [1]_x + [2]_x + [3]_x,$$

$$[1]_x := w \phi \left( x + \frac{3}{2}i \right) \phi \left( x + \frac{1}{2}i \right) Q(x - \frac{3}{2}i) Q(x + \frac{3}{2}i),$$

$$[2]_x := \phi \left( x + \frac{3}{2}i \right) \phi \left( x - \frac{3}{2}i \right) Q(x - \frac{3}{2}i) Q(x + \frac{3}{2}i),$$

$$[3]_x = w^{-1} \phi \left( x - \frac{3}{2}i \right) \phi \left( x - \frac{1}{2}i \right) Q(x + \frac{5}{2}i) Q(x + \frac{1}{2}i).$$

(10)

These boxes are analogous to those for Young tableaux and play a fundamental role in the following.

### 3.1. The symmetric fusion structure

In the bootstrap procedure of the $S$-matrix, the $\phi^3$ property is most essential. A similar feature, referred to as the $a_2^{(2)}$ fusion structure, has been discussed in the dilute $A_L$ model [38]. This comes from the singularity of the RSOS weights at $u = \pm 2$. 

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Due to the fusion procedure, the eigenvalue of a fusion QTM can be represented by sum over products of 'boxes' with different letters and spectral parameters. The point is that the sum of such boxes can be identified with semi-standard Young tableaux (SSTs) for $sl_3$. Let us present a simple example. See [39, 41, 40, 38, 55] for details. By the fusion procedure, one can construct a transfer matrix of which auxiliary space acts on a symmetric subspace of $V \times V$. We associate this with the set of the SSTs, $\begin{bmatrix} i_1 & i_2 \end{bmatrix} (i_1 \leq i_2)$. The eigenvalue of the transfer matrix is then represented by

$$\sum_{i_1 \leq i_2} \begin{bmatrix} i_1 & i_2 \end{bmatrix} x_{i_1-i} x_{i_2+i}.$$ 

In the same manner, one can introduce fusion models based on general Young diagrams, of which eigenvalues are specified by their shapes. On each diagram, the spectral parameter changes by $+2i$ from left to right and $-2i$ from top to the bottom.

First, we argue QTMs associated with rectangular shape diagrams. Although they are simpler, some of the properties below are crucial for the discussion for QTMs associated with complex Young diagrams. First, due to identities,

$$\begin{align*}
\frac{1}{2} x+i &= \phi_2(x) \begin{bmatrix} 1 \end{bmatrix}, \\
\frac{1}{3} x+i &= \phi_2(x) \begin{bmatrix} 2 \end{bmatrix}, \\
\frac{2}{3} x+i &= \phi_2(x) \begin{bmatrix} 3 \end{bmatrix}, \\
\phi_2(x) &= \phi(x \pm \frac{3}{2}i), \\
\frac{1}{2} x+2i &= \phi_3(x) := \prod_{j=1}^{3} \phi(x \pm (\frac{3}{2} - j)i), \\
\frac{2}{3} x-2i &= \phi_3(x) := \prod_{j=1}^{3} \phi(x \pm (\frac{3}{2} - j)i),
\end{align*}$$

the QTMs from a $2 \times m$ ($3 \times m$) Young diagram can be reduced to those from $1 \times m$ (or just scalars). Hereafter, we adopt abbreviations for any function $f$,

$$f(x \pm ia) := f(x + ia)f(x - ia).$$

The first three relations bring additional symmetry besides $sl_3$, and make the $a_2^{(2)}$ symmetry explicit. We call this the $a_2^{(2)}$ structure hereafter.

Second, the eigenvalues of $1 \times m$ fusion QTMs have 'duality' in the following sense. Let us denote a renormalized $1 \times m$ fusion QTM by $T_m(x)$:

$$T_m(x) = \frac{1}{f_m(x)} \sum \begin{bmatrix} i_1 & i_2 & \cdots & i_m \end{bmatrix} x_{i_1-i} x_{i_2+i} \cdots x_{i_m-i}$$

where the semi-standard (SS) condition $i_1 \leq i_2 \leq \cdots \leq i_m$ is imposed on the summation. The spectral parameters are assigned as $x - i(m-1) \cdots x + i(m-1)$ from left to right. The renormalization factor, which is the common factor of the expressions from tableaux of length $m$, is given by

$$f_m(x) := \prod_{j=1}^{m-1} \phi\left(x \pm i \left(\frac{2m-1}{2} - j\right)\right),$$

such that the resultant $T_m$ are all degree $2N$ w.r.t. $h[x + \text{shift}]$. For later convenience, we also denote by $T_{m}^{\text{non}}(x)$, the un-renormalized $T_m(x)$, namely, the object without the factor $f_m(x)$ in (13). Obviously, we have a periodicity due to Boltzmann weights:

$$T_m(x + Pi) = T_m(x) \quad (T_{m}^{\text{non}}(x + Pi) = T_{m}^{\text{non}}(x))$$

where $P$ is defined in (9).
Hereafter, we frequently use functions of which the argument is shifted by a half period in the imaginary direction, i.e., \( x + i\frac{P}{2} \). We shall denote them by a symbol \( \vee \). For example,

\[
T_m^\vee(x) = T_m\left(x + \frac{P}{2}i\right).
\]

The following functional relations are a direct consequence of the \( a_2^{(2)} \) structure.

\[
T_m(x - i)T_m(x + i) = g_m(x)T_m(x) + T_{m+1}(x)T_{m-1}(x), \quad m \geq 1,
\]

\[
g_m(x) = \phi(x \pm i(m + \frac{3}{2})),
\]

\[
T_{-1}(x) := 0,
\]

\[
T_0(x) := f_2(x).
\]

The scalar function \( g_m(x) \) satisfies

\[
g_{2L-1-m}(x) = \begin{cases} g_m(x), & \text{if } L \text{ is even} \\ g_m^\vee(x), & \text{if } L \text{ is odd.} \end{cases}
\]

The adjacency condition leads to \( T_{2L}(x) = T_{2L+1}(x) = 0 \). Thus one deduces the duality relations,

\[
T_m(x) = \begin{cases} T_{2L-1-m}(x), & m = 0, \ldots, 2L \quad \text{for } L \text{ even} \\ T_{2L-1-m}^\vee(x), & m = 0, \ldots, 2L \quad \text{for } L \text{ odd} \end{cases}
\]

for the solutions to equation (14). Note that this duality is not valid for \( T_{\text{non}}(x) \); obviously \( T_{2L-1-m}^{\text{non}}(x) \) and \( T_{m}^{\text{non}}(x) \) have different orders in \( h[x] \). These relations can be in principle proved by using explicit fusion weights. We have at least checked this numerically and assume its validity in this report.

The duality brings about many relations. It thus happens that seemingly quite different equations are identical; an identical object can have (infinitely) many equivalent expressions. This sometimes makes the proof of the functional relations quite involved.

Before closing this section, we make a comment.

The simple relation (14) suggests a choice of \( Y \)-function,

\[
Y_m(x) = \frac{T_{m+1}(x)T_{m-1}(x)}{g_m(x)T_m(x)} \quad (m = 1, \ldots, 2L - 2). \tag{17}
\]

This leads to a set of simple equations (the \( Y \)-system) [55],

\[
Y_m(x - i)Y_m(x + i) = \frac{(1 + Y_{m+1}(x))(1 + Y_{m-1}(x))}{1 + Y_m(x)}. \tag{18}
\]

The relations seem to be tractable and they indeed work well in a nonunitary case [57]. This is not the case for the unitary models. The transformation of the \( Y \)-system into a simple set of integral equations requires both sides to be analytic and non-zero and to have constant asymptotic behaviour as \( |x| \to \infty \) (ANZC) in certain strips.

As is demonstrated numerically in the case of \( L = 8 \) in appendix G, this requirement is not satisfied for (18). Therefore we must seek for alternatives, which may not be found only among symmetric fusion transfer matrices.

In the next subsection, we thus consider a much wider class of skew Young diagrams.

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The resultant table is of the shape $2 \times m$. The spectral parameters are assigned for each box, and the tail (bottom right) are fixed to be $\mu$ from $\lambda$. The spectral parameters of the head (top left) and for the lower diagram. Due to the semi-standard condition, products of $\mu/\lambda$ boxes, as well as the condition $i_k < j_k$ for the upper and $i_1 \leq i_2 \leq \cdots \leq i_m$ and $j_1 \leq j_2 \leq \cdots \leq j_m$, respectively.

For illustration, let us prove (14) diagrammatically. We consider $T^\text{mon}_{m}(x+i)T^\text{mon}_{m}(x-i)$. The spectral parameters are assigned for each box, $x - (m-2)i_1, \ldots, x + mi$ from the left to the right for $T^\text{mon}_{m}(x+i)$, and $x - mi, \ldots, x + (m-2)i$ for $T^\text{mon}_{m}(x-i)$. Thus two diagrams are situated so that these two width $m$ rectangles are piled vertically (figure 2).

We consider a particular assignment of letters, $i_1, \ldots, i_m$ for the upper and $j_1, \ldots, j_m$ for the lower diagram. $T^\text{mon}_{m}(x+i)T^\text{mon}_{m}(x-i)$ is obtained by the sum over tableaux under the condition $i_1 \leq i_2 \leq \cdots \leq i_m$ and $j_1 \leq j_2 \leq \cdots \leq j_m$, respectively.

The first term in the lhs of (14) comes from the configurations $i_k < j_k$, for $k = 1, \ldots, m$. The resultant table is of the shape $2 \times m$ and satisfies the semi-standard condition. Due to
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

The reduction (11), it is a member of $T_{m}^{\text{non}}(x)$ times $\prod_{j=1}^{m} \phi_2(x-i(m+1)+2ij)$. The second term comes from the configurations which break the condition (figure 3). By $i_\alpha \geq j_\alpha$, we mean the first pairs breaking the condition from the left end. Then the ordering within the $i$ and the $j$ concludes that the new sequences $j_1, \ldots, j_\alpha, i_\alpha, i_{\alpha+1}, \ldots, i_m$ and $i_1, \ldots, i_{\alpha-1}, j_{\alpha+1}, \ldots, j_m$ satisfies the semi-standard condition. One can also check that the new sequences satisfy the proper spectral parameter assignment, namely increment of the spectral parameter by $2i$ from left to right. Thus after the recombination of sequences, we obtain $T_{m+1}^{\text{non}}(x)T_{m-1}^{\text{non}}(x)$. Finally, by dividing both sides by the factor $f_m(x \pm i)$, one arrives at (14).

The above procedure is formally interpreted as follows. Place two diagrams properly according to their spectral parameters. We just join them or recombine them such that the total set of spectral parameters is conserved.

The case studies for $L = 3, 4, 6$, however, suggest that $T_{\mu/\lambda}(x)$ is not the right object in the discussion of functional relations. Instead, we should introduce $\Lambda_{\mu/\lambda}(x)$, which is analytic under BAE, from $T_{\mu/\lambda}(x)$ by putting $T_{m \geq 2L}(x) = 0$ in the latter,

$$\Lambda_{\mu/\lambda}(x) := T_{\mu/\lambda}(x)/\{T_{m \geq 2L}(x) = 0\}. \tag{20}$$

The pole-free property of $\Lambda_{\mu/\lambda}(x)$ is obvious from (19).

The diagrammatic decomposition applies to $T_{\mu/\lambda}(x)$ but not to $\Lambda_{\mu/\lambda}(x)$ in a strict sense. Therefore, the decomposition rules, demonstrated above may not seem to be helpful in our discussion. For most of the cases which we treat below, however, the skew diagram $\mu - \lambda$ is not so huge and $\Lambda$ and $T$ coincide. Thus the above rules are quite often applicable and play significant roles; as already remarked, there are plenty of expressions for an identical object. Therefore, it is sometimes crucial to choose ‘the most’ relevant expression among many. The above graphical decomposition provides hints on the optimal candidate in proving the functional relations.

Before closing the section, we mention an important diagrammatic symmetry of $\Lambda_{\mu/\lambda}$.

Corollary 1 Let $Y$ and $\hat{Y}$ be skew diagrams. $\hat{Y}$ is obtained from $Y$ after $180^\circ$ rotation. The following relation is a consequence of equation (19)

$$T_Y(x) = (T_{\hat{Y}}(x^*))^* \tag{21}$$

where the asterisk stands for complex conjugations.

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J. Stat. Mech. (2005) P01004
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

\[
\begin{array}{cccccccccc}
0 & 2 & 4 & 6 & 8 & 10 & 12 & 14 & 16 \\
\end{array}
= \\
\begin{array}{cccccccccc}
0 & 2 & 4 & 6 & 8 & 10 & 12 & 14 & 16 \\
\end{array}
+ \\
\begin{array}{cccccccccc}
2 & 4 & 6 & 8 & 10 & 12 & 14 & 16 & 0 \\
\end{array}
\]

**Figure 4.** The multiplication of $T_1(x)(=T_{B_1})$ by $T_{2L-2}(x+(2L-1)i)$. The letter $a$ in a box implies its spectral parameter $x+ia$.

The proof is direct; one substitutes $\hat{Y}_i = (\mu_1 - \lambda_{d+1-i}) - (\mu_1 - \mu_{d+1-i})$ for $Y = \mu - \lambda$ and $\mu'_1 = d$.

This corollary is crucial in establishing a property of the transfer matrices (see section 8). It will also be used implicitly in the proof of functional relations in the following.

After these preparations, we introduce the first of our three ingredients in the closed functional relations, the breather $T$-system, in the next section.

4. **Breather $T$-system**

The RSOS weights possess a singularity at $u = \pm 3$ besides the $u = \pm 2$ singularity related to $a_2^{(2)}$ structure [38]. It is shown that the singularity at $u = \pm 3$ signifies the ‘hidden’ $sl_2$ structure behind the model. The breather type relation originates from this $sl_2$ structure.

We define the ‘breather’ transfer matrices for the dilute $A_L$ model in regime 2.

\[
T_{B_1}(x) = T_1(x)
\]

\[
T_{B_3}(x) = \begin{cases} 
\frac{1}{\phi(x)}\Lambda_{2L-2,1}\left(x - \frac{5}{2}i\right), & \text{for } L \text{ odd} \\
\frac{1}{\phi(x)}\Lambda_{2L-2,1}\left(x - \frac{5}{2}i\right), & \text{for } L \text{ even}
\end{cases}
\]

\[
T_{B_5}(x) = \frac{1}{\phi(x \pm \frac{3}{2}i)}\Lambda_{(4L-5,2L-2,2L-2)/(2L-3,2L-3)}(x).
\]

One immediately verifies the validity of the following functional relations:

\[
T_{B_1}\left(x + i \frac{L - 2}{2(L + 2)}\right) = T_0\left(x + i \frac{L + 6}{2(L + 2)}\right) + \phi(x)T_{B_3}(x)
\]

(22)

\[
T_{B_3}\left(x + i \frac{L - 2}{2(L + 2)}\right) = T_0(x)T_0\left(x + i \frac{4}{(L + 2)}\right) + T_{B_1}(x)T_{B_5}(x).
\]

(23)

**Proof.** The proof utilizes the duality (16). For (22), we first consider the multiplication of $T_1(x)(=T_{B_1})$ by its dual $T_{2L-2}(x+(2L-1)i)$. See figure 4 for the case $L = 5$.

The spectral parameters are assigned so that they align horizontally. For $L$ even (odd) $T_{2L-2}(x+(2L-1)i)$ is equal to $T_{B_1}(x+(2L-1)i)$ ($T_{B_1}(x+(2L-1)i + \frac{3}{2}i)$). Thus the product is equivalent to that of two $T_{B_1}$. It is decomposed into two terms; the first is $T_{2L-1}$ and the second is $T_{B_3}$ (with some shifts in spectral parameters and with coefficients). By

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Figure 5. The graphical proof of the second relation.

Figure 6. A graphical representation of (26). The letter $a$ in a box implies its spectral parameter $x + ia$.

duality, $T_{2L-1}$ is replaced by $T_0$. Then one arrives at (22) after the shift in the spectral parameter and taking account of proper coefficients.

The second relation is shown similarly. One only should notice that the dual of the diagram $(2L-1,1)$ is given by a skew diagram $(2L-1,2L-1)-(2L-2)$. The case $L=5$ is depicted in figure 5.

Despite its appearance, the first diagram on the rhs corresponds to the products of $T_0$ due to (19). Then it is straightforward to verify (23).

By noticing $3 = -(L-2)(L+2)/(L+1)$, we understand the above relation stems from the singularity of the RSOS weight at $u = \pm 3$ related to the $sl_2$ structure.

The following representation of $T_{B_3}(x)$ and $T_{B_5}(x)$ will be useful in the following argument for regime 2:

$$T_{B_3}\left(x - \frac{L-2}{2(L+2)}i\right) = \frac{1}{\phi(x - \frac{3}{2}i)}\Lambda_{(2,1)}(x + i) \tag{24}$$

$$T_{B_5}(x) = \Lambda_{(3,2)}/(1)(x). \tag{25}$$

Proof. The first equation utilizes the representation of $T_{B_1}(x)$ by the height 2 diagram.

$$T_{B_1}(x)T_{B_1}(x + 3i) = f_2(x + 2i)\phi_2(x)\Lambda_{(2,1)}(x + i) + \frac{\phi_3(x + i)}{\phi_2(x)}. \tag{26}$$

See (11) and (12) for $\phi_{2,3}$.

We present graphically equation (26) in figure 6. By noticing $3 = -(L-2)(L+2)/(L+1)$, and comparing the above with (22) (after a trivial shift in $x$), one concludes (24).
After proper normalizations, it reads where the duality relation (16) is used in the second equality. \( \phi_3(x) \) is defined in (12). After proper normalizations, it reads

\[
T(4L-5,2L-2,2L-2)/(2L-3,2L-3) = T_2(x \pm 2i). \tag{27}
\]

On the other hand, formula (19) leads to

\[
T(4L-5,2L-2,2L-2)/(2L-3,2L-3) = \Lambda_{(4L-5,2L-2,2L-2)/(2L-3,2L-3)} + T_0(x \pm i)T_{4L-3}. \tag{28}
\]

Note that \( T_{4L-3} = T_3(x) \), which is derived in parallel to the duality. Then the comparison (27) and (28) leads to

\[
\Lambda_{(4L-5,2L-2,2L-2)/(2L-3,2L-3)} = T_2(x \pm 2i) - T_0(x \pm i)T_3 = \Lambda_{(3,2)/(1)}(x),
\]

where (19) is applied in the last equality. In this manner, one verifies (25). \( \square \)

The transformation of the \( T \)-system to the ‘breather’ \( Y \)-system is standard [35, 36, 52]. We define

\[
Y_{B_1}(x) := \frac{\phi^\gamma(x)T_{B_1}(x)}{T_0(x \pm i\frac{L+6}{2(L+2)})},
\]

\[
Y_{B_3}(x) := \frac{T_{B_1}(x)T_{B_3}(x)}{T_0(x)T_0(x \pm i\frac{4}{L+2})}.
\]

Then the following ‘breather’ \( Y \)-system is directly derived with the help of the ‘breather’ functional relation (23).

\[
Y_{B_1}(x \pm i\frac{L-2}{2(L+2)}) = (1 + Y_{B_3})(x). \tag{29}
\]

The main goal in this report is to evaluate the largest eigenvalue of \( T_1(u,0) \) or \( T_{B_1}(0) \) in the Trotter limit. For this, our strategy is to utilize a set of closed functional relations, including \( T_{B_1}(x) \). The \( sl_2 \) originated functional relations are indeed enough to obtain the desired relations in the case of \( M_{3,4} + \phi_{1,2} \). This is no longer valid for the general case. In the next sections, we introduce further objects and relations to achieve the goal.

5. The \( sl_3 \) type functional relations

As is mentioned in the preceding section, there is another structure analogous to the \( sl_3 \) models. Let \( W_{m}^{(a)} \), \( (a = 1, 2, m \in \mathbb{Z}_{\geq 0}) \) be a Yangian module corresponding to \( m\Lambda_a \) as a classical \( sl_3 \) module. We associate a transfer matrix \( T_{m}^{(a)}(x) \) with \( W_{m}^{(a)} \). Then the following functional relations hold:

\[
T_{m}^{(a)}(x \pm i\gamma) = T_{m}^{(a+1)}(x)T_{m}^{(a-1)}(x) + T_{m+1}^{(a)}(x)T_{m-1}^{(a)}(x)
\]

where \( \gamma \) is a parameter related to a singularity in the \( R \) matrix. \( T_{m}^{(0)}, T_{m}^{(3)} \) are given by trivial scalar factors multiplied by an identity operator. Thus when multiplied on...
a common eigenspace of the commuting transfer matrices they amount to known trivial scalar functions. The above relation can be graphically represented as follows. Consider a product of two rectangles of the shape $a \times m$. It is decomposed into two subsets. One is the products of two rectangles of the shape $(a + 1) \times m$ and the shape $(a - 1) \times m$ and the other is the products of rectangles $a \times (m + 1)$ and $a \times (m - 1)$.

We turn back to the dilute $A_L$ models in regime 2. The above relations among rectangles cannot hold; due to the $a_2^{(2)}$ symmetry, the $2 \times m$ rectangle is always reduced to the rectangle of $1 \times m$. Nevertheless, an analogue can still be found. Let us denote

$$T_{D_1}(x) = \Lambda_{(j+1,j)/(1,1)}(x).$$

In this notation, which is adopted hereafter, $T_{B_3}(x)$ in the preceding section is identical to $T_{D_2}(x)$. Note also that $T_{D_1}(x) = T_1(x \pm 2i)$ and $T_{D_6}(x) = T_0(x \pm 2i)$.

**Lemma 1** The following functional relations, similar to $sl_3$ type, hold.

$$T_{2j}(x \pm 2i) = T_{2(j+1)}(x)T_{2(j-1)}(x) + T_{D_{2j}}(x)$$  \hspace{1cm} (30)

$$T_{D_{2j}}(x \pm 2i) = \begin{cases} 
T_{D_{2(j+1)}}(x)T_{D_{2(j-1)}}(x) + T_{2j}(x)T_{B_3}(x) \left( x \pm \frac{iL - 2 - 4j}{2L + 4} \right) & \text{j even} \\
T_{D_{2(j+1)}}(x)T_{D_{2(j-1)}}(x) + T_{2j}(x)T_{B_3}^\nu(x) \left( x \pm \frac{iL - 2 - 4j}{2L + 4} \right) & \text{j odd}. 
\end{cases}$$  \hspace{1cm} (31)

The similarity is clearer by writing $T_j^{(1)}(x) = T_{2j}(x)$ and $T_j^{(2)}(x) = T_{D_{2j}}(x)$. The difference lies in the fact that $T_j^{(3)}(x)$, which amounts to a known scalar for the $sl_3$ case, remains non-trivial; it is given by a product of $T_{B_3}$.

**Proof.** The first equality (30) is trivial as shown by the quantum Jacobi–Trudi formula. The second one needs one step further. First, one uses the following relation, proved by the quantum Jacobi–Trudi formula:

$$T_{D_{2j}}(x \pm 2i) - T_{D_{2(j+1)}}(x)T_{D_{2(j-1)}}(x) = T_{2j}(x)T_{(2j+2,2j+1,2j)/(2,1)}(x).$$  \hspace{1cm} (32)

From the tableau rule, a hatched part of the diagram in $T_{(2j+2,2j+1,2j)/(2,1)}(x)$ reduces to a product of scalars; see figure 7.

The resultant two pieces of diagrams represent the products of $\Lambda_{(2,1)}(x + (2j + 1)i)$ and $\Lambda_{(2,3)/(1)}(x - (2j + 1)i)$ apart from normalization factors. Note that $2j \equiv \frac{4j}{2L+4} \pmod{P}$.

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for $j$ even and $2j \equiv \frac{4j}{2L+4} + \frac{P}{2} \pmod{P}$ for $j$ odd. Then from the expression (24) one finds for the first piece

$$\Lambda_{(2,1)}(x + (2j + 1)i) = \phi \left( x + \left( 2j - \frac{5}{2} \right) i \right) T_{B_3} \left( x - \frac{L - 2 - 4j}{2L + 4} i \right).$$

The second piece, $\Lambda_{(2,2)/(1)}(x - (2j + 1)i)$, is equal to $(\Lambda_{(2,1)/(1)}(x + (2j + 1)i))^*$ by corollary 1. Therefore, after taking account of proper normalization factors and using the property $(T_{B_3}(x))^* = T_{B_3}(x^*)$, one finds

$$T_{B_3} \left( x \pm \frac{L - 2 - 4j}{2L + 4} i \right) = \begin{cases} T_{B_3}(2j+2,j+2,j+1)(x) / (2j+1,j+1)(x) & \text{j even}, \\ T_{B_3}(2j+2,j+2,j+1)(x) / (2j+1,j+1)(x) & \text{j odd}. \end{cases}$$

Thus the second relation is proved from (32) and (33).

For $L = 4k - 1, 4k$ and $4k + 1$, we introduce

$$Y_{j}^{(1)}(x) := \begin{cases} \frac{T_{D_{2j}(2j+1)}(x)}{T_{D_{2j}(2j+1)}(x)} & \text{j even} \\ \frac{T_{D_{2j}(2j+1)}(x)}{T_{D_{2j}(2j+1)}(x)} & \text{j odd} \end{cases}$$

$$Y_{j}^{(2)}(x) := \begin{cases} \frac{T_{D_{2j}(2j+1)}(x)T_{B_3}(x \pm \frac{L - 2 - 4j}{2L + 4} \) & \text{j even} \\ \frac{T_{D_{2j}(2j+1)}(x)T_{B_3}(x \pm \frac{L - 2 - 4j}{2L + 4} \) & \text{j odd} \end{cases}$$

while these do not show up for $L = 4k + 2$.

By noticing $2 \equiv \frac{4}{2L+4} + \frac{P}{2}$, one immediately obtains the following $Y$-system:

$$Y_{j}^{(1)} \left( x \pm \frac{4}{2L + 4} i \right) = \frac{1 + Y_{j}^{(2)}(x)}{(1 + \frac{1}{Y_{j}^{(1)}(x)})(1 + \frac{1}{Y_{j}^{(2)}(x)})}$$

$$Y_{j}^{(2)} \left( x \pm \frac{4}{2L + 4} i \right) = \frac{1 + Y_{j}^{(1)}(x)}{(1 + \frac{1}{Y_{j}^{(1)}(x)})(1 + \frac{1}{Y_{j}^{(2)}(x)})} \quad 1 \leq j \leq k - 3.$$

These are desired relations. See the $Y$-system in appendix A.

The $sl_2$ $T$-system is not itself closed. Here we find it connected to the $sl_3$-like $T$-system through $T_{B_3}(x) = T_{D_2}(x)$. The latter system is, however, not closed either. In the next section, other functional relations are introduced, referred to as ‘magnon-like’, which turn out to ‘close’ the total $T$-systems.

6. The magnon-like $t$-system and the Dynkin-like diagrams

The above transfer matrices are not enough to obtain sets of closed functional relations, which are necessary to derive TBA. We have to introduce further objects separately depending on four families. They are related to the magnon-like $Y$-system introduced in [25,30]. A magnon-like $Y$-system is associated with a diagram of the Dynkin type. Although they can be classified into the four families, the diagrams are all different
depending on values of \( L \). Consequently, the magnon-like \( Y \)-systems are all distinct for different \( L \).

Here the situation is slightly unified. The diagrams we have to deal with are ‘tails’ of those in [25,30], independent of \( L \), except for the dilute \( A_{4k+2} \) models.

Before starting discussions on individual cases, we shall make remarks. A blank node in each Dynkin-like diagram indicates its different character from the rest. With node \( a \) in the diagram, we associate \( t_m(a,x) \), where \( m \) takes an integer value from a finite set. Our magnon-like \( t \)-system is a set of functional relations among \( t_m(a,x) \), which should be described separately for four families. The most crucial observation below will be that \( t_m(a) \) associated with a blank node is always expressed by (product of) \( T_B \).

The dilute \( A_{4k-1} \) model and the dilute \( A_{4k+1} \) model share the same diagram but possess different \( t \)-systems. The only ‘ \( L \) ’ dependent diagram for \( A_{4k+2} (L = 4k + 2) \) remarkably coincides with the Dynkin diagram for \( D_{k+2} \).

Below we take \( x \) to be real. As usual, the symbol \( a \sim b \) means that nodes \( a \) and \( b \) are adjacent on the diagram.

6.1. The magnon-like \( t \)-system for \( A_{4k-1} \)

We define a magnon-like \( t \)-system associated with figure 8,

\[
t_m^{(a)}\left(x \pm \frac{i}{2L+4}\right) = t_m^{(a)}(x) t_{m+1}^{(a)}(x) + \prod_{b \sim a} t_m^{(b)}(x), \quad a = 1 \sim 6, \quad m = 1,2 \tag{38}
\]

\[
t_0^{(a)} = 1 \quad (a = 1 \sim 6), \quad t_3^{(a)} = 0, \quad (a = 3 \sim 6)
\]

\[
t_1^{(7)}\left(x \pm \frac{i}{2L+4}\right) = t_2^{(7)}(x). \tag{39}
\]

This is not a set of closed functional relations. However, we note its affinity to the \( E_6 \) related \( T \)-system [36]. Indeed, if \( t_1^{(7)}(x) = 1 \) and \( t_3^{(1)}(x) = t_3^{(2)}(x) = 0 \), the above functional relation coincides with the \( T \)-system for the level 3 RSOS model, corresponding to the coset \((E_6^{(1)})_2 \times (E_6^{(1)})_1/(E_6^{(1)})_3\). Our interpretation is as follows. The non-trivial \( t_1^{(7)}(x) \) makes the \( t_3^{(1)}(x) \) and \( t_3^{(2)}(x) \) non-zero and connects the \( E_6 \) level 3 \( T \)-system to others. Thus \( t_1^{(7)}(x) \) plays the role of ‘glue’. This seems to be quite parallel to the \( sl_3 \)-like \( T \)-system, which is connected to the \( sl_2 \ \( T \)-system by non-trivial \( T_{B_3} \). Here the situation is indeed the same; we will see below that \( t_1^{(7)}(x) = T_{B_3} \). \( T_{B_3} \) thus plays a role in gluing the \( sl_3 \), the \( E_6 \) and the \( sl_2 \) \( T \)-systems.
We shall define ‘gauge invariant’ functions for later discussions [36, 56].

\[
Y_{k-1}^{(a)}(x) := \frac{\prod_{b \sim a} t_{2}^{(b)}(x)}{t_{2}^{(a)}(x)}, \quad a = 1 \sim 6
\]

\[
Y_{k-2}^{(1)}(x) := \frac{t_{2}^{(3)}(x)}{t_{1}^{(1)}(x)t_{2}^{(1)}(x)}
\]

\[
Y_{k-2}^{(2)}(x) := \frac{t_{2}^{(4)}(x)t_{2}^{(7)}(x)}{t_{1}^{(2)}(x)t_{2}^{(2)}(x)}
\]

6.2. The magnon-like t-system for \(A_{4k}\)

Associated with figure 9 we propose a magnon-like t-system for this case,

\[
t_{m}^{(a)} \left( x \pm \frac{i}{L + 2} \right) = t_{m-1}^{(a)}(x)t_{m+1}^{(a)}(x) + \prod_{b \sim a} t_{m}^{(b)}(x), \quad a = 1 \sim 6, \quad m = 1, 2,
\]

\[
t_{0}^{(a)} = 1 \quad (a = 1 \sim 6), \quad t_{3}^{(a)}(x) = 0 \quad (a = 0, 3, 4),
\]

\[
t_{1}^{(7)} \left( x \pm \frac{i}{L + 2} \right) = t_{2}^{(7)}(x).
\]

Again, this is not a set of closed functional relations. In this case we note its similarity to the T-system for the level 3 RSOS model, corresponding to the coset \((A_{5}^{(1)})_{2} \times (A_{5}^{(1)})_{1}/(A_{5}^{(1)})_{3}\) [36]. As in the preceding case, \(t_{1}^{(7)}(x)\) will turn out to be equal to \(T_{B_{3}}(x)\). Thus the interpretation is the same; \(t_{1}^{(7)}\) connects the T-system for the level 3 \(A_{5}\) RSOS model and the \(sl_{3}\) and \(sl_{2}T\)-systems.

Y-functions for this case read

\[
Y_{k-1}^{(a)}(x) := \frac{\prod_{b \sim a} t_{2}^{(b)}(x)}{t_{2}^{(a)}(x)}, \quad a = 1 \sim 6
\]

\[
Y_{k-2}^{(1)}(x) := \frac{t_{2}^{(0)}(x)t_{2}^{(3)}(x)}{t_{1}^{(1)}(x)t_{2}^{(1)}(x)}
\]

\[
Y_{k-2}^{(2)}(x) := \frac{t_{2}^{(0)}(x)t_{2}^{(4)}(x)t_{2}^{(7)}(x)}{t_{1}^{(2)}(x)t_{2}^{(2)}(x)}
\]
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

6.3. The magnon-like $t$-system for $A_{4k+1}$

Although it shares the same diagram with the case $A_{4k-1}$, a magnon-like $t$-system takes a slightly different form,

$$t_1^{(a)} \left( x \pm \frac{i}{2L+4} \right) = t_2^{(a)}(x) + \prod_{a \sim b} t_1^{(b)}(x), \quad a = 3, 4, 5, 6$$

$$t_2^{(1)} \left( x \pm \frac{4i}{2L+4} \right) = t_2^{(1)}(x) + t_1^{(2)}(x)$$

$$t_2^{(2)} \left( x \pm \frac{4i}{2L+4} \right) = t_2^{(2)}(x) + t_1^{(1)}(x)t_1^{(7)}(x)$$

$$t_1^{(7)} \left( x \pm \frac{i}{2L+4} \right) = t_2^{(7)}(x)$$

$$t_2^{(a)} \left( x \pm \frac{i}{2L+4} \right) = \prod_{b \sim a} t_2^{(b)}(x) \quad a = 5, 6$$

$$t_2^{(3)} \left( x \pm \frac{i}{2L+4} \right) = t_1^{(3)}(x)t_2^{(3)}(x) + t_2^{(5)}(x)t_1^{(7)}(x) \left( x \pm \frac{i}{2L+4} \right)$$

$$t_2^{(4)} \left( x \pm \frac{i}{2L+4} \right) = t_1^{(4)}(x)t_2^{(4)}(x) + t_2^{(5)}(x)t_1^{(2)}(x) \left( x \pm \frac{i}{2L+4} \right).$$

It is analogous to the $T$-system for the level 3 RSOS model, corresponding to the coset $(D_4^{(1)})_2 \times (D_4^{(1)})_3$ [36], which is clear by setting $t_1^{(1)} = t_1^{(2)} = t_1^{(7)} = 1$ and $t_3^{(3)} = t_3^{(4)} = 0$. In this case, $t_1^{(7)}$ will turn out to be a product of $T_{B_3}$, while the equations for $t_1^{(1)}$ and $t_1^{(2)}$ will be reminiscent of the underlying $sl_3$ structure. Again, our interpretation is that the non-trivial $T_{B_3}$ glues the $D_4$, the $sl_3$ and the $sl_2$ $T$-systems.

$Y$-functions are defined by

$$Y_{k-1}^{(a)}(x) := \frac{t_2^{(a)}(x)}{\prod_{b \sim a} t_1^{(b)}(x)}, \quad a = 3 \sim 6$$

$$Y_{k-1}^{(1)}(x) := \frac{t_2^{(3)}(x)t_3^{(3)}(x) \left( x \pm \frac{1}{2L+4} i \right)}{t_1^{(5)}(x)t_2^{(1)}(x) \left( x \pm \frac{1}{2L+4} i \right)}$$

$$Y_{k-1}^{(2)}(x) := \frac{t_2^{(4)}(x)t_3^{(4)}(x) \left( x \pm \frac{1}{2L+4} i \right)}{t_2^{(5)}(x)t_1^{(2)}(x) \left( x \pm \frac{1}{2L+4} i \right)}$$

Figure 10. Dynkin-like diagram for $A_{4k+1}$.
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Figure 11. Dynkin-like diagram for $A_{4k+2}$.

$$Y_{k-2}^{(1)}(x) := \frac{t_1^{(2)}(x)}{t_2^{(1)}(x)}$$
$$Y_{k-2}^{(2)}(x) := \frac{t_1^{(1)}(x) t_1^{(7)}(x)}{t_2^{(2)}(x)}.$$

6.4. The magnon-like $t$-system for $A_{4k+2}$

In this case only, the diagram depends on $k$. A magnon-like $t$-system for the figure 11 is defined as

$$t_m^{(a)} \left( x \pm \frac{2i}{L+2} \right) = t_m^{(a)}(x) t_m^{(a)} + \prod_{b \sim a} t_b^{(a)}(x), \quad a = 1 \sim k + 1 \quad m = 1, 2, 3 \quad (43)$$
$$t_0^{(a)} = 1 \quad (a = 1 \sim k + 1), \quad t_4^{(a)}(x) = 0 \quad (a = 2 \sim k + 1)$$
$$t_1^{(0)} \left( x \pm \frac{2i}{L+2} \right) = t_2^{(0)}(x), \quad t_2^{(0)} \left( x \pm \frac{2i}{L+2} \right) = t_1^{(0)}(x) t_3^{(0)}(x). \quad (44)$$

In this case, the functional relation looks like the level 4 RSOS $T$-system for $D_{k+1}$, instead of level 3.

The description of $Y_j^{(a)}$ can be done in the most systematic way among four categories,

$$Y_{a-1}^{(m)}(x) = \frac{t_m^{(a)}(x) t_m^{(a)}(x)}{\prod_{b \sim a} t_b^{(a)}(x)} \quad m = 1, 2 \quad 1 \leq a \leq k + 1.$$

6.5. The solutions to the magnon-like $t$-system

The main finding in this report is that the solutions to the above magnon-like $t$-system are expressible in terms of fusion transfer matrices appearing in dilute $A_L$ models. We summarize the explicit solutions in appendix B. The proof requires several steps and is rather lengthy; part of it will be given in section 10. Before going into such detail, we would like draw attention to a general property of the solution and also its implication for the $Y$-system. Firstly, we need to explain $T_K(x)$ appearing in the solutions for $L = even$.

7. Kink transfer matrix

The case studies on $M_{4,5}, M_{6,7}$ [23] imply the necessity of the introduction of objects which are not produced from boxes by (19). Generally, we find it possible to introduce the desired objects for the $A_L$ model, $L$ even.

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

The explicit procedure to construct is as follows. Let $L = 2\ell$ and consider the model in regime 2, corresponding to the perturbation $\phi_{1,2}$. Take

$$
\omega^\ell \phi(x + \frac{\ell}{\ell + 1} i) \frac{Q(x + \frac{2\ell}{\ell + 1} i)}{Q(x - \frac{2\ell}{\ell + 1} i)}
$$

as the top term (or the highest term). Next, one adds a lower term

$$
\omega^{\ell-1} \phi(x - \frac{\ell + 2}{\ell + 1} i) \frac{Q(x - \frac{2\ell + 1}{\ell + 1} i)Q(x + \frac{2\ell - 1}{\ell + 1} i)}{Q(x - \frac{2\ell}{\ell + 1} i)Q(x + \frac{2\ell - 1}{\ell + 1} i)}
$$

so that the singularity at $x = x_j + \frac{2\ell - 1}{\ell + 1} i$ cancels due to the Bethe ansatz equation (8). Then we introduce another counter-term so as to kill the singularity at $x = x_j - \frac{2\ell - 1}{\ell + 1} i$, and so on.

One observes that this procedure closes by finite steps and a pole-free object, consisting of $2\ell + 1$ terms, results. We refer to this as an eigenvalue of the ‘kink transfer matrix’, $T_K(x)$, although the corresponding transfer matrix is not yet found in the lattice model. For example,

$$
T_K(x) = w^4 \phi(x + \frac{4}{5} i) \frac{Q(x + \frac{2}{5} i)}{Q(x - \frac{2}{5} i)} + w^3 \phi(x - \frac{6}{5} i) \frac{Q(x - \frac{5}{5} i)Q(x - \frac{9}{5} i)}{Q(x - \frac{7}{5} i)Q(x + \frac{4}{5} i)}
$$

for the dilute $A_8$ model.

The numerical investigation on the largest eigenvalue indicates $(Q(x^*))^* = Q(x)$, which leads to an important property $(T_K(x^*))^* = T_K(x)$.

For our argument, it is not important whether it is a ‘real’ transfer matrix or not, but its pole-free property as well as functional relations among the products of kink transfer matrices and $T_m(x)$ are significant. One finds, for instance,

$$
T_K(x + \frac{2k + 1}{4k + 2} i) = 2T_{4k-1}(x)
$$

in the case of the dilute $A_{4k}$ model.

These relations are summarized as lemma 7 in appendix C.

As mentioned there, although the explicit proof has been done only for smaller values of $k$, we assume the following.

**Conjecture 1.** Lemma 7 holds for general $k$.

The functional relations there play an important role in the proof of the magnon-like $t$-system for $L = \text{even}$.

doi:10.1088/1742-5468/2005/01/P01004 24
8. The reality property of magnon-like $t$-functions

We remark the following significant property of our solutions to the $T$-system.

**Lemma 2** Let $T(x)$ be a solution of either the magnon-like $t$-system, the $sl_3$-like $T$-system, or the breather $T$-system. Then, in ‘the largest eigenvalue sector’,

$$T(x) = (T(x^*))^*$$

which we call the reality property. Namely, it assures the reality of $T(x)$ on the real axis.

As is commented in the previous section, $T_K(x)$ satisfies the reality property. One also verifies the property for most of $t^{(a)}_m(x)$ listed in appendix B using the diagrammatic symmetry (21), immediately. It can be shown for the remaining $t^{(a)}_m(x)$, which takes ‘asymmetric’ shape, by utilizing the magnon-like ‘$t$-’ system which can be established without use of the reality property. We take the dilute $A_{4k}$ case for example. Using the diagrammatic symmetry, $t^{(a)}_m(x)$ are shown to satisfy (45) except for $(a, m) = (0, 1)$ and $(a, m) = (4, 1)$. That of the magnon-like $t$-system,

$$t^{(1)}_1\left(x \pm \frac{i}{2L+4}\right) - t^{(1)}_2(x) = t^{(0)}_1(x)t^{(0)}_3(x)$$

can be established independent of the property. Since all other entries other than $t^{(0)}_1(x)$ satisfy (45), it follows from the above relation that $t^{(0)}_1(x)$ should also satisfy (45). The case $(a, m) = (4, 1)$ can be treated similarly. There are, however, two exceptions, $t^{(3)}_1(x)$ and $t^{(4)}_1(x)$ for the dilute $A_{4k+1}$ model. One still can devise other functional relations which explain their reality property. Since the proof is rather technical, we omit it here.

9. The $Y$-system and the $T$-system

Several $Y$-functions are already introduced for four categories by appropriate ratios of $t$- or $T$-functions. Thanks to the $T$-system or the $t$-system, one can easily check that the $Y$-system is satisfied for the case in which the lhs are products of $Y^{(1)}_j$ or $Y^{(2)}_j(x)$ except for $j = k - 2$. One can verify the rest of the $Y$-system by utilizing the explicit expressions of $t$ in appendix B.

Our main message in this report is summarized as follows.

**Proposition 1** The four categories of the $Y$-system proposed in [25] can be solved in terms of transfer matrices appearing in the dilute $A_L$ model, if $x$ is real.

Due to the different forms of $Y$-system, the proof must be given separately for the four categories. Below we shall concentrate on the case $L = 4k$. The other cases can be established with minor modifications.

**Proof.** Proposition 1 for $L = 4k$.

Thanks to explicit relations between $Y$ and $t$ in (40)–(42) the ‘tail’ part of the $Y$-system (A.4) is shown to be a consequence of the magnon-like $t$-system. Equation (A.3)
also follows from the $sl_3$ structure (36) except for $j = k - 2$. Consider the case $\alpha = 1$ and $j = k - 2$. By replacing the $Y$-functions by $t^{(a)}_m$, one finds the lhs,
\[
\frac{t_2^{(0)}(x \pm \frac{2}{L+2} i)}{t_3^{(1)}(x \pm \frac{2}{L+2} i)} t_2^{(3)}(x \pm \frac{2}{L+2} i) t_3^{(1)}(x \pm \frac{2}{L+2} i)
\]
while the rhs reads
\[
\frac{t_2^{(2)}(x \pm \frac{1}{L+2} i)}{t_3^{(2)}(x) t_1^{(1)}(x \pm \frac{2}{L+2} i)} (1 + \frac{1}{Y_{k-3}(x)}).
\]

At this stage, the equality is not obvious. By the explicit forms of $Y_{k-3}^{(1)}(x)$ in (34), we rewrite
\[
1 + \frac{1}{Y_{k-3}^{(1)}(x)} = \begin{cases} \frac{T_{2(k-2)}(x \pm 2i)}{T_{D_{2(k-2)}}(x)}, & k \text{ odd} \\ \frac{T_{2k}^{(\nu)}(x \pm \frac{2\ell+1}{L+2} i) T_{D_{2k}}^{(\nu)}(x \pm \frac{2\ell-3}{L+2} i)}{T_{D_{2k}}^{(\nu)}(x \pm \frac{2\ell+1}{L+2} i) T_{D_{2k}}^{(\nu)}(x \pm \frac{2\ell-3}{L+2} i)}, & k \text{ even.} \end{cases}
\]

By substituting further the explicit forms of $t^{(a)}_m$ in appendix B.2, one verifies the case $a = 1$ and $j = k - 2$; for example, both sides reduce to
\[
\frac{T_{D_{2(k-1)}}(x \pm \frac{2}{L+2} i)}{T_{2(k-2)}(x \pm \frac{2}{L+2} i) T_{2k}(x \pm \frac{2}{L+2} i)}
\]
in the case where $k$ is even. The case $\alpha = 2$ and $j = k - 2$ can be shown in the same manner.

Since equation (A.1) is already proved in (29), it remains to show equation (A.2). We must treat the cases $k$ even and odd separately. For $k = 2m + 1$, the product
\[
\Xi_{k-1}^{(2)}(x) \prod_{\ell=0}^{k-2} \Xi_{k-1}^{(2)}(x \pm \frac{4k - 4\ell - 6}{2L + 4} i) = \Xi_{k-1}^{(2)}(x) \prod_{\ell=1}^{k-1} \Xi_{k-1}^{(2)}(x \pm \frac{4\ell - 2}{2L + 4} i)
\]
is written in terms of fusion transfer matrices using (35),
\[
\frac{T_{D_{2k}}(x \pm \frac{1}{L+2} i)}{T_{D_{2k-2}}^{(\nu)}(x \pm \frac{1}{L+2} i)} \prod_{\ell=1,\ell=\text{odd}}^{k-1} \frac{T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell+1}{L+2} i) T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell-3}{L+2} i)}{T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell+1}{L+2} i) T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell-3}{L+2} i)} \times \prod_{\ell=1,\ell=\text{even}}^{k-1} \frac{T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell+1}{L+2} i) T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell-3}{L+2} i)}{T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell+1}{L+2} i) T_{D_{2k-2}}^{(\nu)}(x \pm \frac{2\ell-3}{L+2} i)}
\]
\[
= \frac{T_{D_{2k}}(x \pm \frac{1}{L+2} i)}{T_{D_{2k-2}}^{(\nu)}(x \pm \frac{1}{L+2} i)} \prod_{\ell=1}^{m} \frac{T_{D_{2k-4\ell+1}}^{(\nu)}(x \pm \frac{4\ell-1-1}{L+2} i) T_{D_{2k-4\ell}}^{(\nu)}(x \pm \frac{4\ell+1}{L+2} i)}{T_{D_{2k-4\ell+1}}^{(\nu)}(x \pm \frac{4\ell-1-1}{L+2} i) T_{D_{2k-4\ell}}^{(\nu)}(x \pm \frac{4\ell+1}{L+2} i)}
\]
\[
= \frac{T_{D_{2k}}(x \pm \frac{2k-1}{L+2} i)}{T_{0}(x \pm \frac{2k-1}{L+2} i)} T_{0}(x \pm \frac{2k}{L+2} i)
\]
where we have used $T_{D_{2k}}^{(\nu)}(x) = T_{0}(x \pm 2i) = T_{0}(x \pm \frac{2}{4k+2} i)$. Multiplying this by the remaining term in the rhs, $(1 + Y_{D_{2k}}(x)) = T_{D_{2k}}(x \pm \frac{2k-1}{L+2} i)/T_{0}(x \pm \frac{2k+3}{L+2} i)$, one reaches the

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26
expression identical to the lhs. Note that $T_{D_2}(x) = T_{B_2}(x)$. Even the case $k$ can be treated analogously. □

Once the solutions to the magnon-like $t$-system are proved to be given by those in appendix B, the solutions to the $Y$-system immediately follow. It is thus vital to verify the solutions. The proofs for the $L = \text{even cases}$ and the odd cases are similar, respectively. We then present the proof for the case $L = 4k + 2$ in the next section, where it can be done most systematically. As a representative of the odd case, we supplement the outline of the proof for the $L = 4k - 1$ case in appendix E.

10. Proof of the magnon-like $t$-system for $L = 4k + 2$

We first prepare a few lemmas.

Lemma 3

$$\Lambda(m,m,1)/(m-1)(x) = T_{D_{m-1}}(x)\phi \left( x \pm \frac{2m + 5}{2}i \right).$$

This lemma holds not only for $L = 4k + 2$, but all $L$. The proof is direct by the $a_2^{(2)}$ property in (11). Although it is not used in this section, this lemma is frequently used in the proof of the magnon-like $t$-system. See appendix D for example.

Lemma 4 The following equality holds by mod $P$:

$$\frac{k}{2k + 2} \equiv \begin{cases} 
2k + 3, & k: \text{even} \\
2k + 3 + \frac{P}{2}, & k: \text{odd}
\end{cases} \quad (46)$$

$$\frac{j}{2k + 2} \equiv \begin{cases} 
2j, & j: \text{even} \\
2j + \frac{P}{2}, & j: \text{odd}.
\end{cases} \quad (47)$$

They can be checked trivially.

Lemma 5

$$(t_2^{(k)}(x))^2 = \begin{cases} 
T_{D_{2k}}^{\nu}(x), & k: \text{even} \\
T_{D_{2k}}(x), & k: \text{odd}.
\end{cases} \quad (48)$$

Lemma 6

$$t_1^{(k)}(x)t_3^{(k)}(x) = \begin{cases} 
\frac{1}{2}(T_{D_{2k+2}}(x) - T_{D_{2k-2}}(x)), & k: \text{even} \\
\frac{1}{2}(T_{D_{2k+2}}^{\nu}(x) - T_{D_{2k-2}}^{\nu}(x)), & k: \text{odd}.
\end{cases} \quad (48)$$

The proofs of the above two lemmas are given in appendix D.

With these preparations, we shall prove the magnon-like $t$-system equation by equation. Consider first the case $m = 1$ in (43). For $a = 1, \ldots, k - 2$, by paying attention to (47) with $j = 1$, one immediately sees that the assertion is actually equivalent to (30).
Thus only the cases $a = k, k - 1$ need proofs. Using $t_1^{(k)}(x) = T_K(x)/2$ and the functional relation (C.2) with $j = k$, we find

$$t_1^{(k)}(x \pm \frac{1}{2k + 2} i) = \begin{cases} T_{2k-2}^\nu(x) + \frac{1}{2}(T_{2k+2}^\nu(x) - T_{2k-2}^\nu(x)) & \text{for } k: \text{even} \\ T_{2k-2}^\nu(x) + \frac{1}{2}(T_{2k+2}^\nu(x) - T_{2k-2}^\nu(x)) & \text{for } k: \text{odd} \end{cases}$$

The first and the second term in the rhs coincide with $t_1^{(k-1)}(x)$ and $t_2^{(k)}(x)$, respectively, which proves the case $a = k$ of (43). The $a = k - 1$ case is less trivial. Substituting $a = k - 1$ in (B.32) and using (30), one finds

$$t_1^{(k-1)}(x \pm \frac{1}{2k + 2} i) = \begin{cases} T_{2k-2}(x \pm 2i) = T_{2k}(x)T_{2k-4}(x) + T_{D_{2k-2}}(x) & \text{for } k: \text{even} \\ T_{2k-2}(x \pm 2i) = T_{2k}(x)T_{2k-4}(x) + T_{D_{2k-2}}(x) & \text{for } k: \text{odd}. \end{cases}$$ (49)

The second term in the rhs is identified with $t_2^{(k-1)}(x)$.

We note

$$t_1^{(k)}(x)t_1^{(k+1)}(x) = (T_K(x)/2)^2 = \begin{cases} T_{2k}(x) & \text{for } k: \text{even} \\ T_{2k}^\nu(x) & \text{for } k: \text{odd} \end{cases}$$

where (C.2) with $j = k + 1$ is used in the secondary equality. Since $T_{2k-4}(x) = t_1^{(k-2)}(x)$, the first term in (49) can be written as $t_1^{(k-2)}(x)t_1^{(k)}(x)t_1^{(k+1)}(x)$. Combining these observations, we find that the rhs of (49) coincides with $t_2^{(k-1)}(x) + t_1^{(k-2)}(x)t_1^{(k)}(x)t_1^{(k+1)}(x)$, which proves (43) for $a = k - 1, m = 1$.

Next, consider the case $m = 2$ in (43). As in the case of $m = 1$, the relation is equivalent to (31) for $a = 1 \sim k - 2$. One only has to pay attention,

$$T_{D_0}^\nu(x) = T_0\left(x \pm \frac{1}{2k + 2} i\right) = t_2^{(0)}(x)$$

for $a = 1$. In the case of $a = k - 1$, the product of $t_2^{(k-1)}(x)$ reads

$$t_2^{(k-1)}(x \pm \frac{1}{2k + 2} i) = \begin{cases} T_{D_{2k-2}}^\nu(x \pm 2i) = T_{D_{2k-2}}^\nu(x)T_{P_3}\left(x \pm \frac{1}{2k + 2} i\right) + T_{D_{2k}}^\nu(x)T_{D_{2k-4}}^\nu(x) & \text{for } k: \text{even} \\ T_{D_{2k-2}}^\nu(x \pm 2i) = T_{D_{2k-2}}(x)T_{P_3}\left(x \pm \frac{1}{2k + 2} i\right) + T_{D_{2k}}(x)T_{D_{2k-4}}(x) & \text{for } k: \text{odd} \end{cases}$$

where (31) is again applied. The first term in the rhs is the product $t_1^{(k-1)}(x)t_3^{(k-1)}(x)$, while the second is $t_2^{(k)}(x)t_2^{(k+1)}(x)t_2^{(k-2)}(x)$. This comes from lemma 5, as $t_2^{(k)}(x)t_2^{(k+1)}(x) = (t_2^{(k)}(x))^2$. Thus the case $a = k - 1, m = 2$ is proved.

The case $a = k, m = 2$, is the most non-trivial. We treat only the case $k$ even for brevity. When taking the product $t_2^{(k)}(x \pm i \frac{1}{2k + 2})$, we substitute two different expressions for $t_2^{(k)}(x)$. The first comes from the already proved relation, $t_2^{(k)}(x) = t_1^{(k)}(x \pm i \frac{1}{2k + 2}) - t_1^{(k-1)}(x)$. The second is as listed in (B.33).
The first expression results, after applying (C.2),

\[ t_2^{(k)} \left( x \pm i \frac{1}{2k+2} \right) = \frac{1}{2} \left( T_{2k}(x)T_{2k-4}(x) + T_{2k}(x)T_{2k+4}(x) - T_{2k-2}(x - \frac{1}{2k+2}i) \right) \times T_{2k+2}^{\nu}(x + \frac{1}{2k+2}i) - T_{2k-2}^{\nu}(x + \frac{1}{2k+2}i) T_{2k+2}^\nu(x - \frac{1}{2k+2}i) \]

= \frac{1}{2} (T_{2k}(x)T_{2k-4}(x) + T_{2k}(x)T_{2k+4}(x) - T_{2k-2}(x - 2i) T_{2k+2}(x + 2i) - T_{2k-2}(x + 2i) T_{2k+2}(x - 2i)).

(50)

From the second expression, \( t_2^{(k)}(x) = (T_{2k+2}(x) - T_{2k-2}(x))/2 \), we find

\[ t_2^{(k)} \left( x \pm i \frac{1}{2k+2} \right) = \frac{1}{2} (T_{2k}(x)T_{2k-4}(x) + T_{2k}(x)T_{2k+4}(x) + T_{D_{2k-2}}(x) + T_{D_{2k+2}}(x) - T_{2k-2}(x - 2i) T_{2k+2}(x + 2i) - T_{2k-2}(x + 2i) T_{2k+2}(x - 2i)) \]

where (30) is used. Doubling (51) and subtracting (50), one arrives at

\[ t_2^{(k)} \left( x \pm i \frac{1}{2k+2} \right) = \frac{1}{2} (T_{D_{2k-2}}(x) + T_{D_{2k+2}}(x)) = T_{D_{2k-2}}(x) + \frac{1}{2} (T_{D_{2k+2}}(x) - T_{D_{2k-2}}(x)). \]

The first term in the rhs coincides with \( t_2^{(k-1)}(x) \), while the second agrees with \( t_1^{(k)}(x)t_3^{(k)}(x) \) due to lemma 6. Thus the case \( a = k, m = 2 \), is established.

Finally consider the case \( m = 3 \). Except for the case \( a = 1 \), the relation is satisfied trivially. When \( a = 1 \), the lhs reads

\[ t_3^{(1)} \left( x \pm \frac{1}{2k+2}i \right) = T_{B_3}(x \pm \frac{k-2}{2k+2}i) T_{B_3}(x \pm \frac{k}{2k+2}i) \]

= \( T_{B_3}(x \pm \frac{k-2}{2k+2}i) T_0(x) T_0(x \pm \frac{2}{2k+2}i) + T_{B_3}(x) T_{B_3}(x) \)

= \( T_{B_3}(x \pm \frac{k-2}{2k+2}i) T_0(x) T_0(x \pm \frac{2}{2k+2}i) \]

+ \( T_{D_2}(x) T_{B_3}(x) T_{B_3}(x \pm \frac{k-2}{2k+2}i) \)

where (23) is used in the second quality and \( T_{D_2}(x) = T_{B_3}(x) \) is used in the last.

By noting \( t_3^{(0)}(x) = T_0(x) T_0(x \pm \frac{2}{2k+2}i) \), one find the resultant expression coincides with \( t_3^{(2)}(x) t_3^{(0)}(x) + t_2^{(1)}(x) t_4^{(1)}(x) \).

We thereby verified the magnon-like \( t \)-system for \( L = 4k+2 \).

11. Thermodynamic Bethe ansatz in the scaling limit

11.1. Analytic properties

The \( Y \)-systems are obtained in [25] by the simply Fourier transformation of the thermodynamic Bethe ansatz equation (TBA, for short), which is proposed in advance. In the present report, the \( Y \)-system has been proved first, in the preceding sections. The (inverse) transformation of the \( Y \)-system to TBA seems thus to be direct. This is, however,
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

not necessary true. One needs further analytic properties of the $Y$-functions as we shall see below.

The $Y$-system typically assumes the following form. (See appendix A for notations.)

$$Y_m^{(a)}(x \pm i\alpha) = \prod_{b,j} \Xi_j^{(b)}(x \pm i\gamma_j^{(b)}) \prod_{c,k} L_k^{(c)}(x \pm i\gamma_k^{(c)}). \tag{52}$$

We try to take the logarithm of both sides and take the Fourier transformation. For the necessary condition that the last transformation can be done simply, one must ensure that the lhs (rhs) is analytic and non-zero in the strip $\text{Im } x \in [-\alpha, \alpha]$ (\text{Im } x \in [-\gamma_j^{(b)}, \gamma_j^{(b)}]). Otherwise, one must take account of contributions from zeros in the strips.

As is already shown, complex combinations of ‘transfer matrices’ solve both sides. The explicit expressions for eigenvalues of transfer matrices thus enable the investigations of the analytic properties of the $Y$-systems, once parameters $N, q$ and $\beta$ are fixed.

We have carried out numerical investigation for ranges of parameters $L = 5, 7, 8, 9, 10, 4 \leq N \leq 10, 0 \leq q \leq 0.5$ and $\beta = 0.04, 0.05, \ldots, 0.1$. Although our main interest is focused on the limit $N \to \infty$, our data suggest weak dependence on $N$. For illustration, appendix G presents some numerical data and graphs showing locations of zeros for various transfer matrices for the dilute $A_8$ model in regime 2. The analytic property of $Y$ can easily be read off from these zeros.

To state our conclusion, we remark that the $Y$-system is invariant, for even $N$, if $Y$ is replaced by $\tilde{Y}$, defined by

$$\tilde{Y}_{B_1}(x) = \begin{cases} 
\frac{Y_{B_1}(x)}{\kappa(x \pm i(1 + u') \frac{L-2}{2(L+2)})} & \text{for } u < 0 \\
Y_{B_1}(x)\kappa\left(x \pm i(1 - u') \frac{L-2}{2(L+2)}\right) & \text{for } u > 0 
\end{cases}$$

and all other cases, $\tilde{Y} = Y$. The parameter $u'$ stands for $\frac{2(L+2)}{L-2} u$. This is due to the definition of $\kappa$,

$$\kappa(x) = \left( \frac{i}{2} \left( i \frac{L+2}{2L-2} \pi x, \tau' \right) \right)^N$$

which satisfies $\kappa(x \pm i \frac{L-2}{2(L+2)}) = 1$. The elliptic nome $q' = \exp(-\tau'), \tau' = \frac{2(L+1)}{L-2} \tau$ is introduced so as to respect the periodicity of the $Y$ function on the real direction of $x$.

We refer the resultant functional relation to the $\tilde{Y}$-system.

Our conjecture is as follows.

**Conjecture 2** Both sides of the $\tilde{Y}$-system are analytic and non-zero for appropriate strips, for any $L, N, q$ and $\beta$.

11.2. The coupled integral equations

Once the conjecture 2 is taken for granted, the transformation of the $\tilde{Y}$-system into a set of integral equations is straightforward. When all $\tilde{Y}$-functions are replaced by $Y$,
only the integral equation for $\log Y_{B_1}(x)$ possesses a drive term. That is, for $u > 0$, with appropriate kernel functions $K_{B_1,a,m}$, it takes the form

$$\log Y_{B_1}(x) = -\log \kappa \left( x \pm i(1-u') \frac{L - 2}{2(L + 2)} \right) + \sum_{(a,m)} \int_{-\tau''}^{\tau''} K_{B_1,a,m}(x - x') \log \Xi_m^{(a)}(x') \, dx'$$

$$+ \sum_{(a,m)} \int_{-\tau''}^{\tau''} K_{B_1,a,m}(x - x') \log L_m^{(a)}(x') \, dx'$$

(53)

where $\tau'' = \frac{(L-2)}{(L+2)} \pi$. For $u < 0$, the driving term should be $\log \kappa(x \pm i(1+u') \frac{L - 2}{2(L + 2)})$.

The free energy is evaluated from

$$\log T_{B_1}(x) = \log r_N(x) + \int_{-\tau''}^{\tau''} s(x - x') \log \varepsilon_N(x') \, dx'$$

$$+ \int_{-\tau''}^{\tau''} s(x - x') \log(1 + Y_{B_1}(x')) \, dx'$$

(54)

$$r_N(x) = \left( \frac{h(x \pm (\frac{L-2}{2L+4} + u)i)h(x \pm (\frac{4L}{2L+4} + u)i)}{h((\frac{L-2}{2L+4})i)h((\frac{4L}{2L+4})i)} \right)^{N/2}$$

$$\varepsilon_N(x) = \left( \frac{h(x \pm (\frac{L-2}{2L+4} - u)i)h(x \pm (\frac{4L}{2L+4} - u)i)}{h((\frac{L-2}{2L+4} + u)i)h((\frac{4L}{2L+4} + u)i)} \right)^{N/2}$$

$$s(k) = \frac{1}{2 \cosh \frac{L-2}{2(L+2)} k}$$

Here we denote the Fourier transformation of $F(x)$ by $\hat{F}(k)$,

$$\hat{F}(k) = \int_{-\tau''}^{\tau''} F(x) e^{-ikx} \, dx$$

$$F(x) = \frac{\delta}{2\pi} \sum_n \hat{F}(k_n) e^{ik_n x}, \quad k_n = n\delta$$

and $\delta = \pi/\tau''$. The free energy is determined by the value of $T_{B_1}(x)$ at the origin, $\beta f = -\log T_{B_1}(0)$, thus we are only interested in $x \sim 0$ in (54).

All the above results are valid for any $N$. Now the Trotter limit can be performed analytically. The drive term in (53) becomes

$$-\log \kappa \left( x \pm i(1-u') \frac{L - 2}{2(L + 2)} \right) \to 4\beta \pi s(x)$$

while

$$\log r_N(0) \to \lambda \beta \left( \log \vartheta_1 \left( \frac{L - 2}{L + 2} \right) \right)' + \left( \log \vartheta_1 \left( \frac{2L}{L + 2} \right) \right)'$$

$$\log \varepsilon_N(x) \to \lambda \beta \left( \log \vartheta_1(\lambda(ix - \frac{L-2}{2L+4})) \right)' + \left( \log \vartheta_1(\lambda(ix + \frac{3L+2}{2L+4})) \right)'$$

for $u = \beta/N$. The crossing parameter $\lambda$ is set to be $\frac{\pi(L+2)}{4(L+1)}$. 

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11.3. Thermodynamic Bethe ansatz equations

We consider the scaling limit \( q \to 0, \beta \to \infty \) [19]. The precise tuning of these two parameters will be specified below. In order to take this limit, we conveniently rewrite the integrals in the preceding subsection in the form ‘kernel function times small functions’. The numerical calculation indicates, in the vicinity of the origin,\[
\{ |Y| \gg 1, \quad u > 0 \}\]
\[
\{ |Y| \ll 1, \quad u < 0 \}
\]
if \( \beta \gg 1 \). In the connection with \( M_{L,L+1} + \Phi_{1,2} \) theory, we are interested in the regime 2 \( (u > 0) \) case. Thus we first rewrite the \( Y \)-system such that the relations involve only \( (1 + \frac{1}{y|_L}) \). Now the important contribution for the integral comes from the ‘Fermi’ surface \( x' \sim \tau'' \). Thus we are interested in the scaling function, \( \lim_{q \to 0} g_{B_i}(x + \tau'') \) etc, and integral equations, among them the TBA equation. To derive these, we take the logarithm of both sides of the transformed \( Y \)-system above, also take the Fourier transform and find
\[
\hat{M} \left( \begin{array}{c}
\log Y_{B_1} \\
\log Y_{B_3} \\
\vdots \\
\log Y_{B_3}
\end{array} \right) = 4\pi \beta \left( \begin{array}{c}
1 \\
0 \\
\vdots \\
0
\end{array} \right) + \hat{K}_0 \left( \begin{array}{c}
\hat{L}_{B_1} \\
\hat{L}_{B_3} \\
\vdots \\
\hat{L}_{B_3}
\end{array} \right) \tag{55}
\]
where \( \hat{L}_{B_1} = \log(1 + \frac{1}{y|_L}) \) and similarly for the others. Note that \( M \) and \( K_0 \) are asymmetric matrices. See appendix H for example for \( L = 5, 7 \). By multiplying \( M^{-1} \) from the left, the kernel matrix for TBA, \( M^{-1}K_0 \) remarkably turns out to be symmetric. This property is crucial in applying the dilogarithm technique to evaluate the central charge. Its elements agree with the expression described in [25] in terms of \( S \)-matrices, under identification \( x = 30/\pi \) in the limit \( q \to 0 \). Finally, we have checked that non-zero drive terms exist only for \( Y \)-functions corresponding to breathers and kinks. For example, we consider \( L = 5 \). The explicit forms of the drive terms in the Fourier space read
\[
\hat{d}_{B_1} = \frac{8\pi \beta \cosh \frac{11}{14} k}{(2 \cosh \frac{11}{14} k - 1)D(k)} \\
\hat{d}_{B_3} = \frac{4\pi \beta (2 \cosh \frac{2}{14} k + 1) (2 \cosh \frac{4}{14} k - 1)}{D(k)} \\
\hat{d}_{B_5} = \frac{16\pi \beta \cosh \frac{11}{14} k \cosh \frac{4}{14} k}{(2 \cosh \frac{2}{14} k - 1)D(k)} \\
\hat{d}_{B_2} = \frac{8\pi \beta \cosh \frac{11}{14} k}{(2 \cosh \frac{2}{14} k - 1)D(k)} \\
\hat{d}_{K_1} = \frac{4\pi \beta}{(2 \cosh \frac{2}{14} k - 1)D(k)} \\
\hat{d}_{K_2} = \frac{8\pi \beta \cosh \frac{4}{14} k}{(2 \cosh \frac{2}{14} k - 1)D(k)},
\]
where we denote by \( \hat{d}_{B_1} \) the drive term associated with \( \log Y_{B_1} \) and so on. A common denominator \( D(k) \) denotes
\[
D(k) = 2 \cosh \frac{12}{14} k + 2 \cosh \frac{10}{14} k - 2 \cosh \frac{6}{14} k - 2 \cosh \frac{4}{14} k + 1.
\]
doi:10.1088/1742-5468/2005/01/P01004
When taking the inverse Fourier transformation, the nearest zero to the real axis of the denominator is relevant in the ‘scaling’ limit, \( q \to 0 \). We have checked for various \( L \) that the zero of the common denominator \( D(k) \) brings the nearest zero, and always \( k = \pm i \frac{\pi}{3} \).

Let us be more explicit for \( L = 5 \). As announced above, we are interested in the integral equation for \( y(x) = \lim_{q \to 0} Y(x + \tau'') \). Thus the associated drive term reads

\[
d^{\text{scaling}}(x) := \lim_{q \to 0} d(x + \tau'') = \lim_{q \to 0} \frac{1}{2\pi} \sum_{m} \int_{-\infty}^{\infty} \exp \left( 2\pi i m \frac{k'}{\delta} + i(x + \tau'')k' \right) \tilde{d}(k') \, dk'
\]

where the Poisson summation formula is applied. As \( q \to 0 \), \( \delta \) and \( \tau'' \) tend to infinity; only the \( m = 0 \), \( m = -1 \) terms contribute,

\[
d^{\text{scaling}}(x) \sim \lim_{q \to 0} \left( \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i(x+\tau'')k'} \tilde{d}(k') \, dk' + \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i(\tau''-x)} \tilde{d}(k') \, dk' \right).
\]

We enclose the complex contour of the integration in the upper half plane for the first term and in the lower half plane for the second. For \( \tau'' \to \infty \), the nearest poles \( k = \pm i \pi/3 \) bring the dominant contributions. We denote

\[
\frac{1}{D(k)} \sim \frac{r}{(k - i \frac{\pi}{3})i}
\]

then

\[
d^{\text{scaling}}(x) = \frac{8\pi r}{2 \cos \frac{\pi}{2\tau} - 1} \exp \left( -\frac{4}{7} \right) \cosh \frac{\pi}{3} x
\]

for example. Let \( R \) be radius of compactification. Then we take \( q \to 0 \) and \( \beta \to \infty \) while \( m_k \) finite; \( m_k R = \frac{8\pi \beta r}{2 \cos \frac{\pi}{2\tau} - 1} q^{\frac{1}{2}} \). This is the precise meaning of our scaling limit. We write \( \frac{\pi}{3} x = \theta \). Immediately, one verifies that all other drive terms also take the form \( m R \cosh \theta \) and their mass ratio agrees with those in [25].

\[
\begin{align*}
m_{B_1} &= 2m_K \cos \frac{11}{42} \pi & m_{B_2} &= 4m_K \cos \frac{11}{42} \pi \cos \frac{3}{42} \pi \\
m_{B_2} &= 4m_K \cos \frac{11}{42} \pi \cos \frac{3}{42} \pi & m_{B_2} &= 2m_K \cos \frac{11}{42} \pi \\
m_{K_2} &= 2m_K \cos \frac{3}{42} \pi.
\end{align*}
\]

We then write \( y_0(\theta) = \exp(\epsilon(\theta)) \), and verify that the TBA is recovered in the scaling limit.

We comment that for general \( L \), the straightforward calculation shows \( m_K \propto q^{\frac{2(L+1)}{4L+2}} \). This is consistent with the scaling relation, \( \xi = \frac{1}{2-2\Delta_{1,2}} \) for \( m \propto q^z \). Indeed comparison of the result leads to the proper scaling dimension, \( \Delta = \frac{L-2}{4L+2} \).

It remains to check the expression for the scaling free energy.

The scaling limit of free energy comes from the contribution from the last term in (54) at \( x = 0 \).

\[
\int_{-\infty}^{\infty} s(x') \log(1 + Y_{B_1}(x')) \, dx' = \int_{-\infty}^{\infty} s(x') \log Y_{B_1}(x') \, dx' + \int_{-\infty}^{\infty} s(x') \log L_{B_1} \, dx'.
\]

We substitute the inverse transformed (55) into the first term in the rhs above to find

\[
\int_{-\infty}^{\infty} s(x') \log(1 + Y_{B_1}(x')) \, dx' = \int_{-\infty}^{\infty} s(x') d_{B_1}(x') \, dx' + \sum_{\ell} \int_{-\infty}^{\infty} K_\ell(x') \log L_{\ell} \, dx'
\]

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

where $\mathcal{K}_\ell$ reads in the Fourier space

$$\hat{\mathcal{K}}_\ell = \hat{s}(\hat{M}^{-1})_{1,j} \hat{D}_{j,\ell}$$

and

$$\hat{D}_{i,j} := \hat{M}_{i,j} + (\hat{K}_0)_{i,j}.$$  

We adopt indices $y_{B_1} = y_1, y_{B_2} = y_2$, and so on.

The scaling free energy $f_s$, the contribution of excitations near the Fermi surface, is thus identified with

$$f_s = -\lim_{q \to 0} \frac{1}{\beta} \sum_{\ell} \int_{-\infty}^\infty \mathcal{K}_\ell(x' + x'') \log \left( 1 + \frac{1}{y_\ell} (x') \right) dx'.$$

To perform the limit, we use a remarkable relation,

$$\hat{\mathcal{K}}_\ell = \hat{d}_\ell$$

which we checked case by case. Then one repeats the same argument in deriving the drive terms in TBA equation, and reaches

$$f_s = -\frac{1}{2\pi} \sum_{\ell} \int_{-\infty}^\infty m_\ell \cosh \theta \log \left( 1 + \frac{1}{y_\ell} \right) d\theta,$$

under tuning $\beta = \sqrt{\frac{2\pi}{3}} R$. Therefore the field theoretical expression of the free energy is also reproduced.

12. Summary and discussions

In this report, we establish relevant functional relations of the dilute $A_L$ models for arbitrary $L$; to analyse the finite temperature property we demonstrate explicitly that TBA equations for $M_{L,L+1} + \phi_{1,2}$, conjectured by Dorey et al [25], are realized in the scaling limit. The crucial idea is to introduce fusion transfer matrices associated with skew Young tableaux and to investigate the functional relations among them.

There are still many open problems. The explicit identification of string solutions would definitely be one of the most important. The complete study on this will shed some light on the way to proceed for TBA in the case of perturbed nonunitary minimal models. We mention the first step in this direction in [57]. The minimal unitary theory perturbed by $\phi_{2,1}$ may be treated within the same framework. One only has to deal with regime 1 instead of regime 2. We already have some results about the ‘exceptional’ cases in terminology of [25]; however, they are omitted in this report for brevity. We hope to complete the programme on the $\phi_{2,1}$ theory for arbitrary $L$ in a subsequent report.

Before closing this paper, we comment on a possible link to fermionic formulae of the Virasoro characters. In [58]–[62], remarkable results are reported on new representations of the Virasoro characters which originate from fermionic particle bases of the Bethe ansatz equations. See also combinatorial arguments in view of TBA [63, 64] and mathematical generalizations based on the combinatorial objects, the rigged configurations [65]–[69]. The fermionic characters obtained in [58]–[64] seem to be connected to $\phi_{1,3}$ perturbation. In [70, 71], they also explicitly showed the different realization of fermionic formulae which are related to $\phi_{2,1}$ or $\phi_{1,5}$ perturbations.
Since the TBA in [25] offers a description of the associated Hilbert space, it might be natural to expect that a new fermionic representation is possible based on the \( \phi_{1,2} \) perturbation. There exist already positive results based on the simplest case, the dilute \( A_3 \) model [72]. It may be natural to expect the generalizations.

Indeed, we check up to \( O(q^{200}) \) that the following formula is valid:

\[
\chi_{5,6}^{1,1} = \sum q^{(1/2)n_0} \prod_{i=1}^{6} (q)_{n_i} \left[ \frac{n_5 + n_6 + n_8}{2} \right] \left[ \frac{n_7}{2} \right],
\]

where

\[
g = \left( \begin{array}{cccccccc}
4, & 6, & 8, & 4, & 2, & 4, & 0, & 0 \\
6, & 12, & 16, & 8, & 4, & 8, & 0, & 0 \\
8, & 16, & 24, & 12, & 6, & 12, & 0, & 0 \\
4, & 8, & 12, & 8, & 4, & 6, & 0, & 0 \\
2, & 4, & 6, & 4, & 3, & 3, & -\frac{1}{2}, & 0 \\
4, & 8, & 12, & 6, & 3, & 7, & -\frac{1}{2}, & 0 \\
0, & 0, & 0, & 0, & -\frac{1}{2}, & -\frac{1}{2}, & 1, & -\frac{1}{2} \\
0, & 0, & 0, & 0, & 0, & 0, & -\frac{1}{2}, & 1
\end{array} \right).
\]

The summation is taken under the condition \( n_7 \leq \frac{n_5 + n_6 + n_8}{2} \in \mathbb{Z}_{\geq 0}, n_8 \leq \frac{n_7}{2} \in \mathbb{Z}_{\geq 0} \). Note that \( g \) has its origin in the kernel functions of the TBA equation.

A similar formula for \( \chi_{7,8}^{1,1} \) is checked up to \( O(q^{100}) \), which may imply the existence of fermionic formulae for the \( L \) odd case in general. We are not yet able to find such a simple expression for the \( L \) even case. We hope to clarify this in a future publication\(^2\).

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**Appendix A. \( Y \)-system**

As obtained in [25], the \( Y \)-system for the minimal unitary theories perturbed by \( \Phi_{1,2} \) falls into four categories. Below we shall summarize their result with slight changes in notations. Instead of their rapidity variable \( \theta \), we use \( x \) as in the text. They are simply related by \( x = \frac{3\theta}{\pi} \). By \( \mathcal{L} \) we mean \( \mathcal{L} = \frac{1}{1 + \frac{1}{x}} \) for a \( Y \)-function. In section 11, an inverse

\(^2\) We have learnt that Warnaar has independently found a polynomial version of these formulas corresponding to both \( \Phi_{1,2} \) and \( \Phi_{2,1} \) perturbations.
function \( L = 1 + \frac{1}{Y} \) is also introduced but it will not be utilized here. We understand indices associated with \( \mathcal{L} \) specify these for \( Y \)-functions, e.g.,

\[
\mathcal{L}^{(\alpha)}_m(x) = \frac{1}{1 + \frac{1}{Y^{(\alpha)}_m(x)}}
\]

e etc. Similarly \( \mathcal{E}^{(\alpha)}_m(x) \) denotes

\[
\mathcal{E}^{(\alpha)}_m(x) = 1 + Y^{(\alpha)}_m(x).
\]

**Appendix A.1. Case \( (L = 4k - 1, k \geq 2) \)**

The \( Y \)-system is

\[
Y_{B_1}\left( x \pm \frac{L - 2}{2L + 4} \right) = \mathcal{E}^{(2)}_{B_3}(x)
\]

\[
Y_{B_3}\left( x \pm \frac{L - 2}{2L + 4} \right) = \mathcal{E}^{(2)}_{B_3}(x) \prod_{\ell=0}^{k-2} \mathcal{E}^{(2)}_{k-1}(x) \prod_{\ell=0}^{k-2} \left( x \pm \frac{4k - 4\ell - 7}{2L + 4} i \right)
\]

\[
Y_{j}^{(\alpha)}\left( x \pm \frac{4}{2L + 4} i \right) = \left( \mathcal{E}^{(2)}_{B_3}(x) \right) \delta_{j,0} \delta_{\alpha,2} \prod_{\ell=0}^{k-2} \left( \mathcal{L}^{(\alpha)}_{\ell}(x) \right) I_{\alpha,j}^{(A_{k-1})}
\]

\[
\times \left[ \mathcal{L}^{(4)}_{k-1}(x) \mathcal{L}^{(6)}_{k-1}(x) \mathcal{L}^{(5)}_{k-1}(x) \prod_{\ell=0}^{k-2} \mathcal{L}^{(5)}_{\ell}(x) \right] \delta_{j,k-2} \delta_{\alpha,1}
\]

\[
\times \left[ \mathcal{L}^{(3)}_{k-1}(x) \mathcal{L}^{(6)}_{k-1}(x) \mathcal{L}^{(5)}_{k-1}(x) \right] \delta_{j,k-2} \delta_{\alpha,2}
\]

(with \( j = 0, \ldots, k - 2; \ \alpha = 1, 2 \) and \( \bar{\alpha} = 3 - \alpha \))

\[
Y_{k-1}^{(\gamma)}\left( x \pm \frac{1}{2L + 4} i \right) = \left[ \mathcal{L}^{(\gamma)}_{k-2}(x) \right] \delta_{1+\gamma} \prod_{\beta=1}^{\gamma} \left[ \mathcal{E}^{(\beta)}_{k-1}(x) \right] \quad (\gamma = 1, \ldots, 6)
\]

where the last product is taken over nodes on the \( E_6 \) Dynkin diagram, which is obtained from figure 10 by deleting node 7. \( I^{(A_{k-1})} \) stands for the incidence matrix for \( A_{k-1} \), where nodes are indexed by 0 to \( k - 2 \).

**Appendix A.2. Case \( (L = 4k, k \geq 2) \)**

The \( Y \)-system is

\[
Y_{B_1}\left( x \pm \frac{L - 2}{2L + 4} \right) = \mathcal{E}^{(2)}_{B_3}(x)
\]

\[
Y_{B_3}\left( x \pm \frac{L - 2}{2L + 4} i \right) = \mathcal{E}^{(2)}_{B_3}(x) \prod_{\ell=0}^{k-2} \mathcal{E}^{(2)}_{k-1}(x) \prod_{\ell=0}^{k-2} \left( x \pm \frac{4k - 4\ell - 6}{2L + 4} i \right)
\]

\[\text{doi:10.1088/1742-5468/2005/01/P01004}\]
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

\[
Y_j(\alpha)\left(x \pm \frac{4}{2L+4}i\right) = \left[\Xi_{B_3}(x)\right]^{\delta_{j0}\delta_{\alpha^2}}\Xi_j^{(\alpha)}(x) \prod_{\ell=0}^{k-2} \left(\mathcal{L}_\ell^{(\alpha)}(x)\right)^{i_{j,\ell}^{(\alpha-1)}} \times \left[\mathcal{L}_{k-1}^{(0)}(x)\mathcal{L}_{k-1}^{(3)}(x)\mathcal{L}_{k-1}^{(1)}(x) \left(x \pm \frac{2}{2L+4}i\right)\right]^{\delta_{j,k-2}\delta_{\alpha1}} \times \left[\mathcal{L}_{k-1}^{(0)}(x)\mathcal{L}_{k-1}^{(4)}(x)\mathcal{L}_{k-1}^{(2)}(x) \left(x \pm \frac{2}{2L+4}i\right)\right]^{\delta_{j,k-2}\delta_{\alpha2}}
\]

(with $j = 0, \ldots, k - 2$; $\alpha = 1, 2$ and $\bar{\alpha} = 3 - \alpha$) \hfill (A.3)

\[
Y_{k-1}(\gamma)\left(x \pm \frac{2}{2L+4}i\right) = \left[\mathcal{L}_{k-2}^{(\gamma)}(x)\right]^{\delta_{\gamma1}+\delta_{\gamma2}} \prod_{\beta\sim\gamma} \left[\Xi_{k-1}^{(\beta)}(x)\right] \hfill (\gamma = 0, \ldots, 4) \hfill (A.4)
\]

where the last product is taken over nodes on the $A_5$ Dynkin diagram, which is obtained from figure 9 by deleting node 7.

**Appendix A.3. Case ($L = 4k + 1, k \geq 2$)**

The $Y$-system is

\[
Y_{B_1}\left(x \pm \frac{(L-2)}{2L+4}i\right) = \Xi_{B_3}(x)
\]

\[
Y_{B_3}\left(x \pm \frac{(L-2)}{2L+4}i\right) = \Xi_{B_1}(x)\Xi_{k-1}^{(2)}(x) \prod_{\ell=0}^{k-2} \Xi_i^{(2)}(x \pm \frac{4k-4\ell-5}{2L+4}i)
\]

\[
Y_j^{(\alpha)}(x \pm \frac{4}{2L+4}i) = \left[\Xi_{B_3}(x)\right]^{\delta_{j0}\delta_{\alpha^2}}\Xi_j^{(\alpha)}(x) \prod_{\ell=0}^{k-2} \left(\mathcal{L}_\ell^{(\alpha)}(x)\right)^{i_{j,\ell}^{(\alpha-1)}} \times \left[\mathcal{L}_{k-1}^{(3)}(x)\mathcal{L}_{k-1}^{(1)}(x) \left(x \pm \frac{1}{2L+4}i\right)\right]^{\delta_{j,k-2}\delta_{\alpha1}} \times \left[\mathcal{L}_{k-1}^{(4)}(x)\mathcal{L}_{k-1}^{(2)}(x) \left(x \pm \frac{1}{2L+4}i\right)\right]^{\delta_{j,k-2}\delta_{\alpha2}}
\]

(with $j = 0, \ldots, k - 2$; $\alpha = 1, 2$ and $\bar{\alpha} = 3 - \alpha$) \hfill (with $j = 0, \ldots, k - 2$; $\alpha = 1, 2$ and $\bar{\alpha} = 3 - \alpha$)

\[
Y_{k-1}^{(1)}\left(x \pm \frac{2}{2L+4}i\right) = \mathcal{L}_{k-2}^{(1)}(x)\Xi_{k-1}^{(4)}(x)\Xi_{k-1}^{(6)}(x) \Xi_{k-1}^{(5)}(x) \left(x \pm \frac{1}{2L+4}i\right) \Xi_{k-1}^{(3)}(x) \left(x \pm \frac{2}{2L+4}i\right)
\]

\[
Y_{k-1}^{(2)}\left(x \pm \frac{3}{2L+4}i\right) = \mathcal{L}_{k-2}^{(2)}(x)\Xi_{k-1}^{(3)}(x)\Xi_{k-1}^{(6)}(x) \Xi_{k-1}^{(5)}(x) \left(x \pm \frac{1}{2L+4}i\right) \Xi_{k-1}^{(4)}(x) \left(x \pm \frac{2}{2L+4}i\right)
\]

\[
Y_{k-1}^{(3)}\left(x \pm \frac{1}{2L+4}i\right) = \Xi_{k-1}^{(1)}(x)\mathcal{L}_{k-1}^{(5)}(x)
\]

\[
Y_{k-1}^{(4)}\left(x \pm \frac{1}{2L+4}i\right) = \Xi_{k-1}^{(2)}(x)\mathcal{L}_{k-1}^{(5)}(x)
\]
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

\[ Y^{(5)}_{k-1}(x \pm \frac{1}{2L+4}i) = \mathcal{L}^{(6)}_{k-1}(x) \mathcal{L}^{(3)}_{k-1}(x) \mathcal{L}^{(4)}_{k-1}(x) \]
\[ Y^{(6)}_{k-1}(x \pm \frac{1}{2L+4}i) = \mathcal{L}^{(5)}_{k-1}(x). \]

**Appendix A.4. Case ($L = 4k + 2, k \geq 2$)**

\[ Y_{B1}(x \pm \frac{L - 2}{2L+4}i) = \Xi_{B3}(x) \]
\[ Y_{B3}(x \pm \frac{L - 2}{2L+4}i) = \Xi_{B1}(x) \Xi_{k}^{(2)}(x) \prod_{\ell=0}^{k-1} \Xi_{\ell}^{(2)}(x + \frac{4k - 4\ell - 4}{2L+4}i) \]
\[ Y^{(\alpha)}_{j}(x \pm \frac{4}{2L+4}i) = \left[ \Xi_{B3}(x) \right]^{\delta_{j,0} \delta_{\alpha,2}} \Xi_{j}^{(\alpha)}(x) \prod_{\ell \sim j} \mathcal{L}_{\ell}^{(\alpha)}(x) \]

(with $j = 0, \ldots, k; \alpha = 1, 2$ and $\bar{\alpha} = 3 - \alpha$)

where the last product is taken over nodes on the $D_{k+1}$ Dynkin diagram, which is obtained from figure 11 by deleting node 0.

**Appendix B. Explicit solutions to magnon-like $t$-system**

Below, sets of explicit solutions to the magnon-like $t$-system are listed. There are, however, many equivalent expressions as remarked in the text. We write just representatives among them. When other expressions are convenient, some remarks will be made.

**Appendix B.1. Solutions for $A_{L=4k-1}$**

\[ t^{(1)}_{1}(x) = \begin{cases} T_{2k}(x), & k \text{ even} \\ T_{2k}^\lor(x), & k \text{ odd} \end{cases} \]
\[ t^{(2)}_{1}(x) = \begin{cases} T_{D_{2k}}(x), & k \text{ even} \\ T_{D_{2k}}^\lor(x), & k \text{ odd} \end{cases} \]
\[ t^{(3)}_{1}(x) = \begin{cases} \Lambda_{(6k-2,2k)/1}(x - \frac{3}{2L+4}i), & k \text{ even} \\ \Lambda_{(6k-2,2k)/1}^\lor(x - \frac{3}{2L+4}i), & k \text{ odd} \end{cases} \]
\[ t^{(4)}_{1}(x) = \begin{cases} \frac{1}{\phi(x + 2ki + \frac{7}{2}i - \frac{i}{2L+4})} \Lambda_{(4k,4k;2k)/(4k-1,1)}(x - 2ki + i - \frac{i}{2L+4}), & k \text{ even} \\ \frac{1}{\phi(x + 2ki + \frac{7}{2}i - \frac{i}{2L+4})} \Lambda_{(4k,4k;2k)/(4k-1,1)}(x - 2ki + i - \frac{i}{2L+4}), & k \text{ odd} \end{cases} \]

\[ \text{doi:10.1088/1742-5468/2005/01/P01004} \]
Figure B.1. Young diagrams associated with $t_1^{(a)}(x)$ for $k = 2$.

$$t_1^{(5)}(x) = \begin{cases} \Lambda_{(10k-4,6k-2,6k-3)/(4k-1,4k-2)}(x), & k \text{ even} \\ \Lambda_{(10k-4,6k-2,6k-3)/(4k-1,4k-2)}(x), & k \text{ odd} \end{cases}$$ \hspace{1cm} (B.5)

$$t_1^{(6)}(x) = \frac{1}{\phi(x - (L + 2)i - \frac{2i}{2L+4})} \Lambda_{(L,1)}(x - \frac{L + 4}{2L + 4}i)$$ \hspace{1cm} (B.6)

$$t_1^{(7)}(x) = T_{B_3}(x).$$ \hspace{1cm} (B.7)

For a guide to the eye, the corresponding diagrams are illustrated in figure B.1 for $k = 2$.

For $t_2^{(a)}$,

$$t_2^{(1)}(x) = \begin{cases} T_{2k-2}^\vee \left( x \pm \frac{3}{2L + 4}i \right), & k \text{ even} \\ T_{2k-2} \left( x \pm \frac{3}{2L + 4}i \right), & k \text{ odd} \end{cases}$$ \hspace{1cm} (B.8)

$$t_2^{(2)}(x) = \begin{cases} T_{D2k-2}^\vee \left( x \pm \frac{3}{2L + 4}i \right), & k \text{ even} \\ T_{D2k-2} \left( x \pm \frac{3}{2L + 4}i \right), & k \text{ odd} \end{cases}$$ \hspace{1cm} (B.9)

$$t_2^{(3)}(x) = \begin{cases} T_{D2k-2}(x)T_{2k-2}^\vee \left( x \pm \frac{2}{2L + 4}i \right), & k \text{ even} \\ T_{D2k-2}^\vee(x)T_{2k-2} \left( x \pm \frac{2}{2L + 4}i \right), & k \text{ odd} \end{cases}$$ \hspace{1cm} (B.10)

$$t_2^{(4)}(x) = \begin{cases} T_{D2k-2} \left( x \pm \frac{2}{2L + 4}i \right)T_{2k-2}^\vee(x), & k \text{ even} \\ T_{D2k-2}^\vee \left( x \pm \frac{2}{2L + 4}i \right)T_{2k-2}(x), & k \text{ odd} \end{cases}$$ \hspace{1cm} (B.11)

$$t_2^{(5)}(x) = t_2^{(6)} \left( x \pm \frac{1}{2L + 4}i \right)$$ \hspace{1cm} (B.12)
Appendix B.2. Solutions for $A_{L=4k}$

$T_1^{(0)}(x) = \begin{cases} T_1^{(3)}(x - \frac{2i}{L + 2}) T_{2k}^{(3)}(x + \frac{i}{L + 2}) - T_1^{(3)}(x + \frac{2i}{L + 2}) T_{2k-2}^{(3)}(x - \frac{i}{L + 2}), & k \text{ even} \\ T_1^{(3)}(x - \frac{2i}{L + 2}) T_{2k}^{(3)}(x + \frac{i}{L + 2}) - T_1^{(3)}(x + \frac{2i}{L + 2}) T_{2k-2}^{(3)}(x - \frac{i}{L + 2}), & k \text{ odd} \end{cases}$

$T_1^{(1)}(x) = \begin{cases} T_{2k}(x), & k \text{ even} \\ T_{2k}^{(3)}(x), & k \text{ odd} \end{cases}$

$T_1^{(2)}(x) = \begin{cases} T_{2k}(x), & k \text{ even} \\ T_{2k}^{(3)}(x) & k \text{ odd} \end{cases}$

$T_1^{(3)}(x) = \frac{T_K(x)}{\sqrt{2}}$

$T_1^{(4)}(x) = \frac{1}{\phi(x - \frac{2L}{L + 2} i)} (T_1^{(3)}(x - \frac{L}{L + 2} i) T_1^{(3)}(x + \frac{i}{2}) - T_1^{(3)}(x + \frac{L + 4}{L + 2} i) T_0^{(3)}(x - \frac{1}{2} i))$

$T_1^{(7)}(x) = T_{B_3}(x)$.

Among $T_2^{(a)}(x)$, the cases $a = 3, 4$ are non-trivial.

$T_2^{(3)}(x) = \begin{cases} T_{2k-2}^{\gamma}(x), & k \text{ even} \\ T_{2k-2}(x), & k \text{ odd} \end{cases}$

$T_2^{(4)}(x) = \begin{cases} T_{2k-2}(x), & k \text{ even} \\ T_{2k-2}^{\gamma}(x), & k \text{ odd} \end{cases}$

The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs and others are easily obtained from them due to the $t$-system,

$$t_2^{(0)}(x) = t_2^{(3)}(x) t_2^{(4)}(x),$$
$$t_2^{(1)}(x) = t_2^{(3)} \left( x \pm i \frac{i}{L+2} \right),$$
$$t_2^{(2)}(x) = t_2^{(4)} \left( x \pm i \frac{i}{L+2} \right).$$

Finally,

$$t_3^{(1)}(x) = \begin{cases} 
(T_{2k-2}^\vee(x))^2 T_{2k-4}(x), & k \text{ even} \\
(T_{2k-2}(x))^2 T_{2k-4}^\vee(x), & k \text{ odd}
\end{cases}$$
$$t_3^{(2)}(x) = \begin{cases} 
(T_{2k-2}(x))^2 T_{2k-4}^\vee(x), & k \text{ even} \\
(T_{2k-2}^\vee(x))^2 T_{2k-4}(x), & k \text{ odd}
\end{cases}$$

Appendix B.3. Solutions for $A_{L=4k+1}$

$$t_1^{(1)}(x) = \begin{cases} 
T_{2k-2}^\vee(x), & k \text{ even} \\
T_{2k-2}(x), & k \text{ odd}
\end{cases}$$
$$t_1^{(2)}(x) = \begin{cases} 
T_{D_{2k-2}}(x), & k \text{ even} \\
T_{D_{2k-2}}^\vee(x), & k \text{ odd}
\end{cases}$$
$$t_1^{(3)}(x) = \begin{cases} 
\Lambda_{(6k+2,2k)/(1)}(x - \frac{1}{2L+4} \left( x + \frac{i}{2} \right)), & k \text{ even} \\
\Lambda_{(6k+2,2k)/(1)}^\vee(x - \frac{1}{2L+4} \left( x + \frac{i}{2} \right)), & k \text{ odd}
\end{cases}$$
$$t_1^{(4)}(x) = \frac{1}{\phi^\vee(x + (4k+3)i)} \Lambda_{(4k,4k,2k)/(4k-1,1)}^\vee \left( x + \frac{i}{2} \right),$$
$$t_1^{(5)}(x) = \begin{cases} 
\Lambda_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x), & k \text{ even} \\
\Lambda_{(10k+2,6k+2,6k+1)/(4k+1,4k)}^\vee(x), & k \text{ odd}
\end{cases}$$
$$t_1^{(6)}(x) = \frac{1}{\phi(x - (L + \frac{3i}{2})i + \frac{3L}{2L+4}i)} \Lambda_{(L,1)} \left( x + \frac{3L}{2L+4}i \right)$$
$$t_1^{(7)}(x) = T_{B_{3}} \left( x \pm \frac{3i}{2L+4} \right).$$

The corresponding diagrams are illustrated in figure B.2 for $k = 2$. Notice their similarity to those in figure B.1.
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Figure B.2. Young diagrams associated with $t_1^{(a)}(x)$ for $k = 2$.

For $t_2^{(a)}$,

\begin{align*}
  t_2^{(1)}(x) &= \begin{cases} 
    T_{2k}(x)T_{2(k-2)}(x), & k \text{ even} \\
    T_{2k}^\vee(x)T_{2(k-2)}^\vee(x), & k \text{ odd}
  \end{cases} \\
  t_2^{(2)}(x) &= \begin{cases} 
    T_{D_{2k}}^\vee(x)T_{D_{2(k-2)}}^\vee(x), & k \text{ even} \\
    T_{D_{2k}}(x)T_{D_{2(k-2)}}(x), & k \text{ odd}
  \end{cases} \\
  t_2^{(3)}(x) &= \begin{cases} 
    T_{D_{2k}}^\vee(x)T_{2k}\left(x \pm \frac{2}{2L+4}i\right), & k \text{ even} \\
    T_{D_{2k}}(x)T_{2k}\left(x \pm \frac{2}{2L+4}i\right), & k \text{ odd}
  \end{cases} \\
  t_2^{(4)}(x) &= \begin{cases} 
    T_{D_{2k}}^\vee\left(x \pm \frac{2}{2L+4}i\right)T_{2k}(x), & k \text{ even} \\
    T_{D_{2k}}\left(x \pm \frac{2}{2L+4}i\right)T_{2k}^\vee(x), & k \text{ odd}
  \end{cases} \\
  t_2^{(5)}(x) &= t_2^{(6)}\left(x \pm \frac{1}{2L+4}i\right) \\
  t_2^{(6)}(x) &= \begin{cases} 
    T_{D_{2k}}^\vee(x)T_{2k}(x), & k \text{ even} \\
    T_{D_{2k}}(x)T_{2k}^\vee(x), & k \text{ odd}
  \end{cases}
\end{align*}

and the rest are found to be

\begin{align*}
  t_3^{(3)}(x) &= t_2^{(5)}(x) \\
  t_3^{(4)}(x) &= t_2^{(5)}(x)T_{B_3}(x).
\end{align*}
Appendix B.4. Solutions for $A_{L=4k+2}$

In this case, the eigenvalues of $t$ associated with the tails, $k$ and $k+1$, are identical.

\[
t^{(k)}_1(x) = t^{(k+1)}_1(x) = \frac{T_K(x)}{2}\]
\[
t^{(a)}_1 = \begin{cases} 
T_{2a}(x), & a \text{ even} \\
T_{2a}^\vee(x), & a \text{ odd}
\end{cases}  \quad 0 \leq a \leq k - 1
\]

\[
t^{(k)}_2(x) = t^{(k+1)}_2(x) = \begin{cases} 
\frac{T_{2k+2}(x) - T_{2k-2}(x)}{2}, & k \text{ even} \\
\frac{T_{2k+2}(x) - T_{2k-2}(x)}{2}, & k \text{ odd}
\end{cases}
\]

\[
t^{(a)}_2 = \begin{cases} 
T_{D_{2a}}(x), & a \text{ even} \\
T_{D_{2a}}(x), & a \text{ odd}
\end{cases} \quad 0 \leq a \leq k - 1
\]

\[
t^{(k)}_3(x) = t^{(k+1)}_3(x) = T_{B_3}(x)
\]

\[
t^{(a)}_3(x) = T_{B_3}\left(x \pm \frac{k-a}{2k+2}i\right), \quad 1 \leq a \leq k - 1
\]

\[
t^{(1)}_4(x) = T_{B_1}(x)T_{B_3}\left(x \pm \frac{k-2}{2k+2}i\right).
\]

Appendix C. Functional relations for the kink transfer matrix

We verify the following lemma for relatively small values of $k$.

**Lemma 7** The following functional relations among kink transfer matrices and $T_m$ are valid for the dilute $A_{4k}$ and $A_{4k+2}$ models ($k = 1, 2, 3$ for the former and $k = 0 \sim 3$ for the latter) in regime 2 (corresponding to $M_{L,L+1} + \phi_{1,2}$).

**$A_{4k}$:**

\[
T_K\left(x \pm \frac{2k+1}{4k+2}i\right) = 2T_{3k-1}(x)
\]

\[
T_K\left(x \pm \frac{2(k-j)+1}{4k+2}i\right) = \begin{cases} 
2(T_{2j-2}(x) + T_{4k-2j}^\vee(x)), & j: \text{odd} \\
2(T_{2j-2}(x) + T_{4k-2j}(x)), & j: \text{even},
\end{cases} \quad j = 1, \ldots, k
\]

\[
T_K\left(x \pm \frac{2(k+j)+1}{4k+2}i\right) = \begin{cases} 
2(T_{2j-1}(x) + T_{4k-2j-1}^\vee(x)), & j: \text{odd} \\
2(T_{2j-1}(x) + T_{4k-2j-1}(x)), & j: \text{even}
\end{cases} \quad j = 1, \ldots, k.
\]
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

$A_{4k+2}$:

$$T_K\left(x \pm \frac{k + 1}{2k + 2} i\right) = 2T_{2k+1}^\vee(x)$$  \hspace{1cm} (C.1)

$$T_K\left(x \pm \frac{k + 1 - j}{2k + 2} i\right) = \begin{cases} 2(T_{2j-2}(x) + T_{4k-2j+2}(x)), & j: \text{odd}, \quad j = 1, \ldots, k + 1 \\ 2(T_{2j-2}(x) + T_{4k-2j+2}^\vee(x)), & j: \text{even}, \quad j = 1, \ldots, k + 1 \end{cases}$$  \hspace{1cm} (C.2)

$$T_K\left(x \pm \frac{k + 1 + j}{2k + 2} i\right) = \begin{cases} 2(T_{2j-1}(x) + T_{4k-2j+1}(x)), & j: \text{odd}, \quad j = 1, \ldots, k + 1 \\ 2(T_{2j-1}^\vee(x) + T_{4k-2j+1}^\vee(x)), & j: \text{even}, \quad j = 1, \ldots, k + 1 \end{cases}$$  \hspace{1cm} (C.3)

We also note another functional relation for $T_K(x \pm i)$ of the $A_{4k+2}$ case, indicating the $\phi^3$ property

$$T_K(x \pm i) = \begin{cases} 2(\phi^\vee(x)T_K(x) + 2T_{2k-1}^\vee(x)), & k: \text{odd} \\ 2(\phi(x)T_K(x) + 2T_{2k-1}(x)), & k: \text{even}. \end{cases}$$  \hspace{1cm} (C.4)

There are also relations obtained from the above by using the duality and the periodicity, omitted for brevity.

**Proof.** The proof is given by comparing the explicit dressed vacuum forms on both sides. In most cases, this procedure, although simple enough in appearance, is not straightforward. For example, for the dilute $A_{10}$ model, the product of two $T_K$ with identical spectral parameters decomposes as

$$T_K^2(x) = 3T_4(x) + T_{15}(x).$$

By the duality, $T_{15}(x) = T_4(x)$, one concludes (C.2), with $k = 2$ and $j = 3$. Similarly, consider

$$T_K^2\left(x \pm \frac{i}{6}\right) = 2T_2\left(x + \frac{11}{6} i\right) + T_6\left(x + \frac{11}{6} i\right) + T_{13}\left(x + \frac{11}{6} i\right).$$

The rhs is identical to $2(T_2(x + \frac{11}{6} i) + T_6(x + \frac{11}{6} i))$ due to the duality, which proves (C.2), with $k = 2$ and $j = 2$. The rest are also shown by using this kind of trick. \hfill \Box

**Corollary 2** For $A_{4k+2}$ models, the kink transfer matrix is represented as

$$T_K(x) = \begin{cases} \frac{1}{\phi^\vee(x)}(T_{2k+1}^\vee(x) - T_{2k-1}(x)), & k: \text{odd} \\ \frac{1}{\phi^\vee(x)}(T_{2k+1}(x) - T_{2k-1}(x)), & k: \text{even}. \end{cases}$$  \hspace{1cm} (C.5)

This is shown by comparing (C.3) with $j = k + 1$ and (C.4).
For the dilute $A_L$ models with odd $L$, we can also construct $T_K(x)$ following the above manner. The resulting pole-free object consists of $2L + 2$ terms. However, it proves to be identically zero under some assumption on analyticity. At present, we do not know the implication of this.

Appendix D. Proofs of lemmas 5 and 6

**Proof.** Lemma 5.

We use one of the magnon-like $t$-systems

$$t_1^{(k)}\left(x \pm i \frac{1}{2k+2}\right) = t_2^{(k)}(x) + t_1^{(k-1)}$$

which has been shown in the main text. Instead of using the solution for $t_2^{(k)}(x)$ in (B.33), we solve the above for $t_2^{(k)}(x)$ and take its square,

$$(t_2^{(k)}(x))^2 = \left(t_1^{(k)}\left(x \pm i \frac{1}{2k+2}\right)\right)^2 - 2t_1^{(k)}\left(x \pm i \frac{1}{2k+2}\right)t_1^{(k-1)}(x) + \left(t_1^{(k-1)}(x)\right)^2$$

$$= \frac{1}{4}\left(T_k\left(x + i \frac{1}{2k+2}\right)\right)^2 - \frac{1}{4}\left(T_k\left(x - i \frac{1}{2k+2}\right)\right)^2$$

$$- \frac{1}{2}T_k\left(x \pm i \frac{1}{2k+2}\right)t_1^{(k-1)}(x) + \left(t_1^{(k-1)}(x)\right)^2$$

where the solutions for $t_1^{(k)}$ are substituted. By further substituting the solutions for $t_1^{(k-1)}$ and making use of the functional relations for $T_k$ (the cases $j = k$ and $j = k + 1$ in equation (C.2)), one arrives at

$$(t_2^{(k)}(x))^2 = \begin{cases} 
T_{2k}\left(x \pm i \frac{1}{2k+2}\right) - T_{2(k+1)}(x)T_{2(k-1)}(x), & k: \text{ even} \\
T_{2k}\left(x \pm \frac{1}{2k+2}\right) - T_{2(k+1)}(x)T_{2(k-1)}(x), & k: \text{ odd}.
\end{cases}$$

By noting the mod $P$ relation (47) with $j = 1$ and using the $sl_3$ type relation (30), one concludes $(t_2^{(k)}(x))^2 = T_{D2k}(x)$ for $k$ even and $(t_2^{(k)}(x))^2 = T_{D2k}(x)$ for $k$ odd, which proves lemma 5.

**Proof.** Lemma 6.

Thanks to the $\phi^3$ property (C.5), $t_1^{(k)}(x)$ has an expression

$$t_1^{(k)}(x) = \frac{T_K(x)}{2} = \begin{cases} 
\frac{1}{2\phi^\vee(x)}(T_{2k+1}(x) - T_{2k-1}(x)), & k: \text{ even} \\
\frac{1}{2\phi^\vee(x)}(T_{2k+1}(x) - T_{2k-1}(x)), & k: \text{ odd}.
\end{cases}$$
We solve (22) for $t_3^{(k)}(x) = T_{B_3}(x)$ and substitute the result into the product,

$$t_1^{(k)}(x)t_3^{(k)}(x) = \begin{cases} 
\frac{1}{2(\phi^\vee(x))^2} T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x) - T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k-1}(x) \\
-T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x) + T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k-1}(x) & k: \text{even} \\
\frac{1}{2(\phi^\vee(x))^2} T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x) - T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k-1}(x) \\
-T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x) + T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k-1}(x) & k: \text{odd}. 
\end{cases}$$

(D.1)

We will show that the above expression coincides with the rhs of (48). For simplicity, only the case $k$ even is considered below. Consider the second term in the rhs of (D.1) first. Noting the expression $T_1(x + \frac{k}{2k+2} i) = T_1(x + 2ki)$ due to (47) with $m = k$, we have

$$T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k-1}(x) = T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x) + \Lambda((2k-1)/(2k-2)) \left( x + \frac{k}{2k+2} i \right)$$

$$+ \frac{f_{2k}(x-i)}{f_{2k-1}(x)} \Lambda((2k-2)/(2k-1)) \left( x + \frac{k}{2k+2} i \right) + \frac{f_{2k}(x+i)}{f_{2k-1}(x)} \Lambda((2k-1)) \left( x + \frac{k}{2k+2} i \right).$$

(D.2)

The graphical representation in figure D.1 will be intuitively helpful. The second term in the rhs of (D.2) is identical to $(\phi^\vee(x))^2 T_{D_{2k-2}}(x)$ due to lemma 3 with $m = 2k - 1$. The third and the fourth term in the rhs of (D.2) can be further decomposed with the aid of the quantum Jacobi–Trudi formula and mod $P$ relations in lemma 4.

$$T_0(x - (2k - 1)i)T_{2k}(x - i)T_1(x + 2ki) - T_0(x + (2k - 1)i)T_{2k+1}(x)$$

$$= T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k}(x - i) T_1(x - (2k + 3)i) - T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x)$$

for the third, and

$$T_0(x + (2k + 1)i)T_{2k}(x + i)T_1(x - 2ki) - T_0(x + (2k - 1)i)T_{2k+1}(x)$$

$$= T_0 \left( x - \frac{k}{2k+2} i \right) T_{2k}(x + i) T_1(x + (2k + 3)i) - T_0 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x)$$

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

The graphical representation of equation (D.3) for $k = 2$. A figure $a$ in a box specifies its spectral parameter $x + ia$.

for the fourth in (D.2). Next we consider the first term in the rhs of (D.1). To obtain a non-trivial decomposition, we note the following expression for $T_{2k+1}(x)$ due to the $a_2^{(2)}$ property:

$$T_{2k+1}(x) = \frac{1}{\phi(x \pm \frac{4k+5}{2}i)} \Lambda_{2k+1,2k+1}(x).$$

We also use the expression $T_1(x \pm \frac{k}{2k+2}i) = T_1(x \pm (2k+3)i)$ due to (46). By these tricks, three rectangles are now set in the right position so that non-trivial decomposition occurs.

$$T_1(x \pm \frac{k}{2k+2}i) T_{2k+1}(x) = \phi\left(x \pm \frac{4k+3}{2}i\right) T_{D_{2k+2}}(x) + T_0\left(x \pm \frac{k+2}{2k+2}i\right) T_{2k-1}(x) + n_c \Lambda_{2k+1,2k+1}^{(2k+1)}(x) + n_0 \Lambda_{2k+1,2k+1}^{(2k+1)/(1)}(x) \tag{D.3}$$

The scalar function of the first term in the rhs can also be written as $\phi(x \pm \frac{4k+3}{2}i) = (\phi'(x))^2$.

The graphical representation in figure D.2 will again be helpful. The normalization factor $n_c$ in the third and the fourth term in (D.3) reads

$$n_c = \frac{\phi_3(x - i(2k + 1))}{f_{2k+1}(x) \prod_{j=1}^{2k+1} \phi_2(x - 2(k + 1)i + 2ji)}.$$

The upper index ‘non’ in the third and the fourth term indicates the lack of their normalization factors. See (11) for $\phi_2$. The third term in (D.3) is further transformed with the aid of the relation

$$\Lambda_{2k+1,2k}^{(2k+1)}(x) = \Lambda_{2k+1,2k}^{(2k+2)}(x + i) T_1(x + (2k + 3)i) - \Lambda_{2k+1,2k+1}^{(2k+1)}(x + 3i)$$

$$= f_{2k}(x + i) \prod_{j=1}^{2k} \phi_2(x - 2ki + 2ji) T_1(x + (2k + 3)i) T_{2k}(x + i)$$

$$- f_{2k-1}(x) \phi_3(x + (2k + 1)i) \prod_{j=1}^{2k-1} \phi_2(x - 2ki + 2ji) T_{2k-1}(x)$$

where we have used the tableau rule in the second equality. After taking account of normalization factors and mod $P$ relations in lemma 4, the third term in (D.3) is given by

$$T_0\left(x - \frac{k + 2}{2k + 2}i\right) T_{2k}(x + i) T_1(x + (2k + 3)i) - T_0\left(x \pm \frac{k + 2}{2k + 2}i\right) T_{2k-1}(x).$$
Similarly the fourth term reads
\[ T_0 \left( x + \frac{k + 2}{2k+2} i \right) T_{2k}(x - i)T_1(x - (2k + 3)i) - T_0 \left( x + \frac{k + 2}{2k+2} i \right) T_{2k-1}(x). \]

Therefore, the first two terms in the rhs of (D.1) add up to
\[ T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k+1}(x) - T_1 \left( x + \frac{k}{2k+2} i \right) T_{2k-1}(x) \]
\[ = (\phi'(x))^2 (T_{D_{2k+2}}(x) - T_{D_{2k-2}}(x)) + T_0 \left( x + \frac{k + 2}{2k+2} i \right) (T_{2k+1}(x) - T_{2k-1}(x)). \]

Obviously, the last two terms in the above cancel the third and the fourth term in (D.1). Thus we arrive at
\[ t_1^{(k)}(x)t_3^{(k)}(x) = \frac{1}{2}(T_{D_{2k+2}}(x) - T_{D_{2k-2}}(x)). \]
The case for odd \( k \) is done in parallel. \( \square \)

Appendix E. Proof of the magnon-like \( t \) system for \( L = 4k + 1 \)

We summarize the proofs for the magnon-like \( t \)-system, \( L = 4k + 1 \). Most of the relations in (38), of which the lhs takes the form \( t_2^{(a)}(x \pm \frac{1}{2L+4} i) \), are shown without difficulty; thus they are omitted except for \( a = 3, 4 \).

Note the useful mod \( P \) relations
\[ \frac{-1}{4k+3} \equiv (4k + 1) + \frac{P}{2} \quad \text{(E.1)} \]
\[ \frac{1}{4k+3} \equiv (4k + 3) \quad \text{(E.2)} \]
\[ \frac{4}{2L+4} = \frac{2}{4k+3} \equiv 2 - \frac{P}{2} \quad \text{(E.3)} \]

which will be used frequently below.

Appendix E.1. The proof for \( t_1^{(1)}(x \pm \frac{4}{2L+4} i) = t_1^{(2)}(x) + t_2^{(1)}(x) \)

Thanks to (E.3), the lhs is equal to \( T_{2(k-1)}(x \pm 2i) \) or \( T_{2(k-1)}^{(\nu)}(x \pm 2i) \), depending on \( k \) even or odd. Then the \( sl_3 \) relation (30) proves the assertion.

Appendix E.2. The proof for \( t_1^{(2)}(x \pm \frac{4}{2L+4} i) = t_1^{(1)}(x)t_1^{(7)}(x) + t_2^{(2)}(x) \)

The proof is similar. One only has to apply (31).

Appendix E.3. The proof for \( t_1^{(3)}(x \pm \frac{1}{2L+4} i) = t_1^{(1)}(x)t_1^{(5)}(x) + t_2^{(3)}(x) \)

We consider the case \( k \) odd. Using the realness of \( t_1^{(3)}(x) \), the mod \( P \) relation (E.1) and the diagrammatic symmetry, we rewrite the lhs as
\[ \Lambda_{(6k-2,2k)}(x + (4k + 1)i)\Lambda_{(6k,2k+1)/4k+2}(x - (4k + 1)i), \]

\[ \text{doi:10.1088/1742-5468/2005/01/P01004} \]
while, thanks to the duality, the second term in the rhs is given by

\[ t_2^{(3)}(x) = T_{D_{2k}}(x)T_{6k+1}(x \pm (4k + 3)i). \]

Then the quantum Jacobi–Trudi formula leads to

\[
t_1^{(3)}(x + \frac{1}{2L + 4}i) - t_2^{(3)}(x) = T_{2(k-1)}(x)(T_{2k+2}(x)T_{6k+1}(x \pm (4k + 3)i) + T_{2k-2}(x)T_{6k+3}(x \pm (4k + 1)i)
\]

\[
- T_{2k}(x - 2i)T_{6k+1}(x + (4k + 3)i)T_{6k+3}(x - (4k + 1)i)
\]

\[
- T_{2k}(x + 2i)T_{6k+1}(x - (4k + 3)i)T_{6k+3}(x + (4k + 1)i))
\]

\[= T_{2(k-1)}(x)\Lambda_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x).\]

One can check that the rhs of the above equation is equal to \(t_1^{(1)}(x)t_1^{(5)}(x)\).

**Appendix E.4. The proof for** \(t_2^{(3)}(x \pm \frac{1}{2L + 4}i) = t_1^{(1)}(x \pm \frac{i}{2L + 4})t_2^{(5)}(x) + t_1^{(3)}(x)t_3^{(3)}(x)\)

Again we treat the case \(k\) odd. Since \(t_2^{(5)}(x) = t_3^{(3)}(x)\), the equality to prove is

\[
t_2^{(3)}(x \pm \frac{1}{2L + 4}i) = t_3^{(3)}(x)\left( t_1^{(3)}(x) + t_1^{(1)}(x \pm \frac{i}{2L + 4}) \right). \tag{E.4}\]

Substitute (B.25) and use (B.29) for \(t_3^{(3)}(x)\); the lhs is given by

\[
t_2^{(3)}(x) = t_3^{(3)}(x)T_{2k}(x + (4k + 1)i)T_{2k}(x + 2i) \bigg|_{x \rightarrow x - i/(2L + 4)}
\]

\[= t_3^{(3)}(x)\left( T_{2k}(x + (4k + 1)i)T_{6k+1}(x + (8k + 6)i) \right) \bigg|_{x \rightarrow x - i/(2L + 4)}
\]

where the duality is used in the second equality. By (19), it is readily shown that the rhs of the bracket above coincides with

\[
t_3^{(3)}(x)\left( \Lambda_{(6k+2,2k)/(1)}^{(5)}(x - \frac{i}{2L + 4}) + T_{2k-2}(x \pm \frac{i}{2L + 4}) \right)
\]

It is immediate to check that the above equation coincides with the rhs of (E.4).

**Appendix E.5. The proof for** \(t_1^{(6)}(x \pm \frac{1}{2L + 4}i) = t_1^{(5)}(x) + t_2^{(6)}(x)\)

The first term in the rhs, \(t_1^{(5)}(x)\), is identified with \(\Lambda_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x)\) for \(k\) odd or \(\Lambda_{(10k+2,6k+2,6k+1)/(4k+1,4k)}^{(5)}(x)\) for \(k\) even. We consider instead \(T_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x)\). See (20) for their difference. Due to the tableau rule, it clearly decomposes into three parts.
The dilute $\Lambda_l$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Figure E.1. The decomposition of $T_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x)$ into three parts. The hatched part reduces to a product of $\phi_3(x)$.

The ‘height-three’ part reduces to the shifted product of $\phi_3(x)$ in (12). The remaining two parts are identified with the diagram for $t^{(6)}_1(x)$ and its 180° rotation, referring to (B.21) (see figure E.1). By the reality property of $t^{(6)}_1(x)$ and the diagrammatic symmetry (21), one finds

$$T_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x) = \begin{cases} t^{(6)}_1(x \pm \frac{1}{2L+4}i) & k: \text{odd} \\ (t^{(6)}_1)^\vee(x \pm \frac{1}{2L+4}i) & k: \text{even} \end{cases}$$

where the mod relation has been applied. On the other hand, the quantum Jacobi–Trudi formula concludes

$$T_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x) = \Lambda_{(10k+2,6k+2,6k+1)/(4k+1,4k)}(x) = T_{10k+4}(x)T_{D_{2k}}(x) = T_{2k}^\vee(x)T_{D_{2k}}(x),$$

which coincides with $t^{(6)}_2(x)(t^{(6)}_2)^\vee(x)$ for $k$ odd (even) in (B.28). Thus by comparing the above due expressions, one concludes

$$t^{(6)}_1(x \pm \frac{1}{2L+4}i) = t^{(5)}_1(x) + t^{(6)}_2(x) \tag{E.5}$$

irrespective of whether $k$ is even or odd.

Appendix E.6. The proof for $t^{(4)}_1(x \pm \frac{1}{2L+4}i) = t^{(2)}_1(x)t^{(5)}_1(x) + t^{(4)}_2(x)$

We treat the case $k$ even. First note the following representation of $t^{(4)}_1(x + \frac{1}{2L+4}i)$:

$$t^{(4)}_1(x + \frac{1}{2L+4}i)\phi(x + (6k + \frac{9}{2})i) = \Lambda_{(4k-1,4k-1)/(4k-2)}(x + 2i)T_{2k}(x - (2k + 2)i)$$

$$- \Lambda_{(4k+1,4k+1)/(4k)}(x)T_{2k-2}(x - 2ki).$$

This relation is natural by its diagrammatic representation, and can be verified easily by (19). Again we use the reality property of $t^{(4)}_1(x)$ and the diagrammatic symmetry (21) to evaluate $t^{(4)}_1(x \pm \frac{1}{2L+4}i)\phi(x \pm (6k + \frac{9}{2})i)$. 

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After recomposing the products of terms, it is found that
\[
\begin{align*}
t_1^{(4)}(x \pm \frac{1}{2L+4} i) & \phi(x \pm (6k + \frac{9}{2} i)) = \{T_{2k-2}(x \pm 2i)\Lambda_{(6k+1,6k+1,4k+1,1)/(6k,2k)}(x) \\
& + (T_{2k}(x))^2\Lambda_{(6k+1,6k+1,4k-1,1)/(6k,2k+2)}(x) \\
& - T_{2k}(x)T_{2k-2}(x \pm 2i)\Lambda_{(6k+1,6k+1,4k-1,1)/(6k,2k)}(x) \\
& - T_{2k}(x)T_{2k-2}(x \pm 2i)\Lambda_{(6k+1,6k+1,4k+1,1)/(6k,2k+2)}(x) \\
& + \{T_{2k-4}(x)T_{2k}(x) - T_{2k-2}(x \pm 2i)\}T_{2k}(x)\Lambda_{(6k+2,6k+2,1)/(6k+1)}(x) \} \\
& \{T_{2k}(x)\Lambda_{(6k+1,6k+1,4k-1,1)/(6k,2k)}(x) \\
& + T_{D_{2k-2}}(x)\Lambda_{(L,1)}(x - (2k + 2)i)\Lambda_{(L,L)/(L-1)}(x + (2k + 2)i).}
\end{align*}
\]

(6)

Thanks to (19), the content of the first curly bracket is equal to
\[
T_{2k}(x)\Lambda_{(6k+1,6k+1,4k-1,1)/(6k,2k+2)}(x)
+ T_{D_{2k-2}}(x)\Lambda_{(L,1)}(x - (2k + 2)i)\Lambda_{(L,L)/(L-1)}(x + (2k + 2)i). \tag{6.7}
\]

The first term in the above is decomposed into three parts due to the tableau rule as in the case of \(t_1^{(6)}(x \pm \frac{1}{2L+4} i)\). An analogous argument leads to its expression,
\[
\frac{1}{\phi(x \pm (2k - \frac{9}{2})i)}T_{2k}(x)\Lambda_{(2k+1,2k+1,1)/(2k)}(x \pm i(4k + 1))
= \phi(x \pm (6k + \frac{9}{2} i))T_{2k}(x)T_{D_{2k}}(x \pm (4k + 1)i) = \phi(x \pm (6k + \frac{9}{2} i))t_2^{(4)}(x)
\]

where lemma 3 and the mod \(P\) relation are used in the second equality.

The second term in (E.7) is identified with
\[
\phi(x \pm i(6k + \frac{9}{2} i))(t_2^{(4)}(x) + T_{D_{2k-2}}(x)(t_1^{(5)}(x) + t_2^{(6)}(x)))
\]

by the reality property of \(t_1^{(6)}(x)\) and (21). Substituting the established relation, \(t_1^{(6)}(x \pm \frac{1}{2L+4} i) = t_2^{(6)}(x) + t_1^{(5)}(x)\), we obtain the following expression for the first curly bracket term in (E.6),
\[
\phi(x \pm i(6k + \frac{9}{2} i))(t_2^{(4)}(x) + T_{D_{2k-2}}(x)(t_1^{(5)}(x) + t_2^{(6)}(x)))
\]

\[
= \phi(x \pm i(6k + \frac{9}{2} i))(t_2^{(4)}(x) + t_1^{(5)}(x) + t_1^{(2)}(x)t_2^{(6)}(x)).
\]

The second curly bracket term in (E.6) is easily found to be
\[
T_{D_{2k-2}}(x)T_{2k}(x)\Lambda_{(6k+1,6k+1,4k+1,1)/(6k,2k)}(x) = -\phi(x \pm i(6k + \frac{9}{2} i))T_{D_{2k-2}}(x)T_{2k}(x)T_{D_{2k}}(x)
\]

\[
= -\phi(x \pm i(6k + \frac{9}{2} i))t_1^{(2)}(x)t_2^{(6)}(x).
\]

Lemma 3 is applied first, then the definitions in (B.17) and (B.28) are utilized. Combining these results, we obtain
\[
t_1^{(4)}(x \pm \frac{1}{2L+4} i) = t_1^{(2)}(x)t_2^{(5)}(x) + t_1^{(4)}(x)t_2^{(4)}(x).
\]

Appendix E.7. The proof for \(t_2^{(4)}(x \pm \frac{1}{2L+4} i) = t_1^{(2)}(x \pm \frac{1}{2L+4} i)t_2^{(5)}(x) + t_1^{(4)}(x)t_3^{(4)}(x)\)

Here we set \(k = \text{odd}\). Since \(t_3^{(4)}(x) = T_{B_3}(x)t_2^{(5)}(x)\), the assertion is equivalent to
\[
t_2^{(4)}(x \pm \frac{1}{2L+4} i) = t_2^{(5)}(x)(t_1^{(2)}(x \pm \frac{1}{2L+4} i) + t_1^{(4)}(x)T_{B_3}(x).
\]

To prove this, we prepare a few facts.
Finally, by noticing Lemma 10:

\[
\frac{1}{\phi(x - (m + \frac{3}{2})i)} \Lambda_{(m,m-1)}(x) = \frac{1}{\phi(x + (m + \frac{5}{2})i)} \Lambda_{(m,m)/(m-1)}(x + i)
\]

\[
\frac{1}{\phi(x - (m + \frac{3}{2})i)} \Lambda_{(m,1)}(x) = \frac{1}{\phi(x + (m + \frac{5}{2})i)} \Lambda_{(m,m)/(1)}(x + i).
\]

Both of them are easily proved by the \(a_2^{(2)}\) property.

**Lemma 9**

\[
T_{D_{k+1}}(x)T_{D_2}(x + (4k + 5)i) - T_{D_{k+3}}(x + 2i)T_{D_2}(x + (4k + 3)i)
= \frac{\Lambda_{(2,1)}(x + (6k + 4)i)}{\phi(x + (2k + \frac{11}{2})i)\phi(x + (6k + \frac{5}{2})i)} D
\]

where

\[
D := \left( T_{6k+1}(x + 2i)\Lambda_{(4k+3,4k+3)/(4k+2)}(x - (2k - 1)i) - T_{6k+3}(x + 2i)\Lambda_{(4k+1,4k+1)/(4k)}(x - (2k - 3)i) \right).
\]

It is easily checked that the lhs is equal to

\[
T_{6k+1}(x + 2i)T_{(6k+3,6k+3,6k+1)/(4k+3,4k+1)}(x + 2i) - T_{6k+3}(x)T_{(6k+1,6k,6k-1)/(4k+1,4k)}(x + 4i).
\]

The diagram decomposition of this then leads to the rhs of (E.8).

**Lemma 10**

\[
D = \frac{\phi(x + (2k + \frac{11}{2})i)}{\phi(x - (6k + \frac{5}{2})i)} \Lambda_{(10k,6k+2,6k+1)/(2k,2k)}(x).
\]

We use the two decompositions of \(T_{(10k,6k+2,6k+1)/(2k,2k)}(x)\). The comparison of these two leads to

\[
\Lambda_{(10k,6k+2,6k+1)/(2k,2k)}(x) = \frac{\phi(x - (6k + \frac{1}{2})i)}{\phi(x + (2k + \frac{11}{2})i)} T_{2k}(x - (8k + 2)i)\Lambda_{(4k-1,1)}(x + (6k + 2)i)
- T_{6k+3}\Lambda_{(4k+1,4k)}(x + (-2k + 2)i).
\]

By the duality, the quantum Jacobi–Trudi formula and lemma 8, we find

\[
\Lambda_{(4k-1,1)}(x + (6k + 2)i) = \frac{\phi(x + (2k + \frac{3}{2})i)}{\phi(x + (2k + \frac{5}{2})i)} \Lambda_{(4k+3,4k+3)/(4k+2)}(x - (2k - 1)i)
\]

\[
\Lambda_{(4k+1,4k)}(x + (-2k + 2)i) = \frac{\phi(x - (6k + \frac{1}{2})i)}{\phi(x + (2k + \frac{11}{2})i)} \Lambda_{(4k+1,4k+1)/(4k+2)}(x - (2k - 3)i).
\]

Finally, by noticing \(T_{2k}(x - (8k + 2)i) = T_{6k+1}(x + 2i)\), we prove lemma 10.

The final lemma asserts a relation,

**Lemma 11**

\[
\Lambda_{(10k,6k+2,6k+1)/(2k,2k)}(x - (6k + 5)i) = \Lambda_{(4k,4k,2k)/(4k-1,1)}(x)
= \phi(x + (4k + \frac{5}{2})i) \left( t_1^{(4)} \right)^{\nu}(x - \frac{i}{2}).
\]

doi:10.1088/1742-5468/2005/01/P01004
The first equality is due to duality and the second comes from the definition. Now we prove the relation. We conveniently consider

\[
\left( t_2^{(4)} \left( x \pm \frac{i}{2L+4} \right) - t_1^{(2)} \left( x \pm \frac{i}{2L+4} \right) t_2^{(5)} (x) \right) \bigg|_{x \to x + 3i/(2L+4)}
\]

and show that it is equal to \( t_1^{(4)} (x + \frac{3i}{2L+4}) t_3^{(4)} (x + \frac{3i}{2L+4}) \). First, we rewrite the difference as

\[
t_2^{(5)} \left( x + \frac{3i}{2L+4} \right) T_{D_{6k+1}} (x) T_{D_{2k}} (x + (4k + 5)i) - T_{D_{6k+3}} (x + 2i) T_{D_{2k-2}} (x + (4k + 3)i)
\]

(E.9)

then apply lemma 9. With the expression of \( T_{B_3} (x) \) in (24), and the relation \( t_3^{(4)} (x) = t_2^{(5)} (x) T_{B_3} (x) \), one concludes that (E.9) coincides with

\[
\phi (x + 2k + \frac{11}{2}i) \bigg| D.
\]

Therefore, the relation is equivalent to

\[
\frac{D}{\phi (x + (2k + \frac{11}{2})i)} = t_1^{(4)} \left( x + \frac{3i}{2L+4} \right).
\]

(E.10)

We will show this by rewriting the lhs of (E.10), thanks to lemma 10,

\[
\frac{1}{\phi (x - (6k + \frac{1}{2})i)} \Lambda_{(10k,6k+2k,6k+1)/(2k,2k)} (x).
\]

This can be further rewritten as, with the aid of lemma 11,

\[
\phi (x + (10k + \frac{15}{2})i) \bigg| (t_1^{(4)})^\vee (x + (6k + 5)i - \frac{1}{2})
\]

which is readily shown to be equal to the rhs of (E.10) for \( k \) odd. The case \( k \) even can be treated similarly.

**Appendix E.8. The proof for** \( t_1^{(5)} (x \pm \frac{1}{2L+4}i) = t_1^{(3)} t_1^{(4)} t_1^{(6)} + t_2^{(5)} (x) \)

We assume \( k \) even. The proof for the case \( k \) odd is similar. One substitutes \( t_1^{(5)} (x) = t_1^{(6)} (x \pm \frac{1}{2L+4}i) - t_2^{(6)} (x) \) and uses the relation \( t_2^{(5)} (x) = t_2^{(6)} (x \pm \frac{1}{2L+4}i) \).

Then the assertion is transformed into an equivalent form,

\[
t_1^{(3)} t_1^{(4)} (x) = t_1^{(6)} (x + \frac{2}{2L+4}i) - t_2^{(6)} \left( x + \frac{1}{2L+4}i \right) t_1^{(6)} \left( x - \frac{2}{2L+4}i \right)
\]

\[
- t_1^{(6)} \left( x - \frac{1}{2L+4}i \right) t_1^{(6)} \left( x + \frac{2}{2L+4}i \right).
\]
Proof. Lemma 12.

We first note a decomposition of $t_1^{(4)}(x)$,

\[
\begin{align*}
t_1^{(4)}(x + \frac{1}{2L+4}i) &= \frac{1}{\phi(x + (6k + \frac{9}{2})i)} T_{2k}(x) \Lambda_{(4k-1,4k-1)/(4k-2)}(x + (2k + 4)i) \\
&\quad - T_{2k-2}(x + 2i) t_1^{(6)}(x - \frac{1}{2L+4}i)
\end{align*}
\]

where we have used the diagrammatic symmetry,

\[
t_1^{(6)}(x - \frac{1}{2L+4}i) = \left(t_1^{(6)}(x + \frac{1}{2L+4}i)\right)\ast
\]

\[
= \frac{1}{\phi(x + (6k + \frac{9}{2})i)} \Lambda_{(4k+1,4k+1)/(4k)}(x + (2k + 2)i).
\]

Then the product in the lhs of (E.12) is rewritten as

\[
\begin{align*}
t_1^{(3)}(x + \frac{1}{2L+4}i) t_1^{(4)}(x + \frac{1}{2L+4}i) \\
&= \frac{1}{\phi(x + (6k + \frac{9}{2})i)} T_{2k}(x) \Lambda_{(4k-1,4k-1)/(4k-2)}(x + (2k + 4)i) t_1^{(3)}(x + \frac{1}{2L+4}i) \\
&\quad - T_{6k+3}(x - (8k + 2)i) t_1^{(3)}(x + \frac{1}{2L+4}i) t_1^{(6)}(x - \frac{1}{2L+4}i)
\end{align*}
\]

where the duality relation $T_{2k-2}(x + 2i) = T_{6k+3}(x - (8k + 2)i)$ is applied. Note that $t_1^{(3)}(x + \frac{1}{2L+4}i) \equiv \Lambda_{(6k+2,2k)/(1)}(x)$. Then we apply the decomposition rule to the second
Lemma 13.

Proof. The first equality follows from the quantum Jacobi–Trudi formula. Then the duality and the established relation in appendix E.5 lead to the last line. By substituting this into (E.14), one verifies lemma 12.

From lemma 12 and (E.11), one only has to show

\[ R = t_1^{(6)} \left( x + i \frac{3}{2L + 4} \right) \lambda_{2k}^{(6)}(x). \]  

(E.15)

Our argument consists of two further steps. First we will show a lemma.

Lemma 13

\[
R = \frac{T_{2k}(x)}{\phi(x - (2k + \frac{3}{2})i)\phi(x - (6k + \frac{1}{2})i)\phi(x + (6k + \frac{3}{2})i)}
\times \Lambda_{(L,1)}(x + (2k + 1)i) \left( T_1(x - (10k + 5)i)\Lambda_{(2k,2k)/(2k-1)}(x - (8k + 2)i) \right.
- \left. T_0(x - (10k + 6)i)\Lambda_{(2k-1,2k-1)/(2k-2)}(x - (8k + 1)i) \right). \]  

(E.16)

Then the rhs of the above will be shown to be equal to that of (E.15).

Proof. Lemma 13.

The definition of \( R \) in (E.13) is not convenient in applying decomposition rules. We utilize the duality and rewrite

\[
t_1^{(6)} \left( x - \frac{1}{2L + 4} \right) \phi \left( x + \left( 6k + \frac{9}{2} \right) i \right) = \Lambda_{(2k,4k)}(x - (2k + 2)i)
\Lambda_{(4k-1,4k-1)/(4k-2)}(x + (2k + 4)i) = \Lambda_{(2k,4k+2)}(x - (2k + 2)i).
\]

Then they are located at right positions to apply the quantum Jacobi–Trudi formula. The resultant expression reads

\[
R = \frac{T_{2k}(x)}{\phi(x + (6k + \frac{3}{2})i)\phi(x - (2k + 1)i)\phi(x - (6k + \frac{1}{2})i)}
\times \left( T_{8k}(x - (2k + 1)i)\Lambda_{(8k,4k+1,4k+1,4k)}/(2k,2k-1)(x - 2ki) \right.
- \left. T_{8k+1}(x - (2k + 2)i)\Lambda_{(8k,4k,4k-1)}/(2k-1,2k-2)(x + (-2k + 1)i) \right)
\times \left( T_{1}(x - (10k + 5)i)\Lambda_{(8k,4k+1,4k+1,4k)}/(2k,2k-1)(x - 2ki) \right.
- \left. T_{0}(x - (10k + 6)i)\Lambda_{(8k,4k,4k-1)}/(2k-1,2k-2)(x + (-2k + 1)i) \right).
\]

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

\begin{equation}
\begin{array}{c}
4k+1 \\
\downarrow \\
2k \\
\downarrow \\
2k \\
\end{array}
\end{equation}

\begin{align*}
R &= T_{2k}(x)\Lambda_{(4k+1,1)}(x+(2k+1)i) \\
&\quad \times \left( T_1(x-(10k+5)i)\Lambda_{(2k,2k)/(2k-1)}(x-(8k+2)i) \\
&\quad - T_0(x-(10k+6)i)\Lambda_{(2k-1,2k-1)/(2k-2)}(x-(8k+1)i) \right).
\end{align*}

Thus the lemma is established. \hfill \Box

The final step starts from rewriting the first term in the bracket of (E.16) in lemma 13 as
\begin{align*}
\frac{\Lambda_{(1,1)}(x-(10k+5)i)}{\phi_2(x-(10k+5)i)} \Lambda_{(2k,2k)/(2k-1)}(x-(8k+2)i) \\
&= \frac{\phi(x-(10k+\frac{7}{2})i)}{\phi(x-(10k+\frac{15}{2})i)} \Lambda_{(2k+1,2k+1,1)/(2k)}(x-(8k+4)i) \\
&\quad + T_0(x-(10k+6)i)\Lambda_{(2k-1,2k-1)/(2k-2)}(x-(8k+1)i).
\end{align*}

The second term in the rhs coincides with that in the bracket of (E.16). Thus
\begin{align*}
R &= \frac{T_{2k}(x)t_{1}^{(6)}(x+\frac{3i}{2k+4})}{\phi(x-(6k+\frac{1}{2})i)\phi(x-(10k+\frac{15}{2})i)} \Lambda_{(2k+1,2k+1,1)/(2k)}(x-(8k+4)i)
\end{align*}

where $\phi(x-(10k+\frac{7}{2})i) = \phi(x+(6k+\frac{9}{2})i)$ and the definition of $t_1^{(6)}$ is used. Finally, we apply lemma 3 to the above and reach expression (E.15). Thereby we prove the functional relation.

\text{where the duality is applied in the last equality.}

The ‘height 3’ diagrams are decomposed into smaller pieces due to the tableau rule; see figure E.2.

Accordingly, $R$ is represented by
\begin{align*}
R &= \frac{T_{2k}(x)\Lambda_{(4k+1,1)}(x+(2k+1)i)}{\phi(x+(6k+\frac{9}{2})i)\phi(x-(2k+\frac{3}{2})i)\phi(x-(6k+\frac{1}{2})i)} \\
&\quad \times \left( T_1(x-(10k+5)i)\Lambda_{(2k,2k)/(2k-1)}(x-(8k+2)i) \\
&\quad - T_0(x-(10k+6)i)\Lambda_{(2k-1,2k-1)/(2k-2)}(x-(8k+1)i) \right).
\end{align*}

The second term in the rhs coincides with that in the bracket of (E.16). Thus
\begin{align*}
R &= \frac{T_{2k}(x)t_{1}^{(6)}(x+\frac{3i}{2k+4})}{\phi(x-(6k+\frac{1}{2})i)\phi(x-(10k+\frac{15}{2})i)} \Lambda_{(2k+1,2k+1,1)/(2k)}(x-(8k+4)i)
\end{align*}

where $\phi(x-(10k+\frac{7}{2})i) = \phi(x+(6k+\frac{9}{2})i)$ and the definition of $t_{1}^{(6)}$ is used. Finally, we apply lemma 3 to the above and reach expression (E.15). Thereby we prove the functional relation.

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Accordingly, $R$ is represented by
\begin{align*}
R &= \frac{T_{2k}(x)\Lambda_{(4k+1,1)}(x+(2k+1)i)}{\phi(x+(6k+\frac{9}{2})i)\phi(x-(2k+\frac{3}{2})i)\phi(x-(6k+\frac{1}{2})i)} \\
&\quad \times \left( T_1(x-(10k+5)i)\Lambda_{(2k,2k)/(2k-1)}(x-(8k+2)i) \\
&\quad - T_0(x-(10k+6)i)\Lambda_{(2k-1,2k-1)/(2k-2)}(x-(8k+1)i) \right).
\end{align*}

The second term in the rhs coincides with that in the bracket of (E.16). Thus
\begin{align*}
R &= \frac{T_{2k}(x)t_{1}^{(6)}(x+\frac{3i}{2k+4})}{\phi(x-(6k+\frac{1}{2})i)\phi(x-(10k+\frac{15}{2})i)} \Lambda_{(2k+1,2k+1,1)/(2k)}(x-(8k+4)i)
\end{align*}

where $\phi(x-(10k+\frac{7}{2})i) = \phi(x+(6k+\frac{9}{2})i)$ and the definition of $t_{1}^{(6)}$ is used. Finally, we apply lemma 3 to the above and reach expression (E.15). Thereby we prove the functional relation.

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The ‘height 3’ diagrams are decomposed into smaller pieces due to the tableau rule; see figure E.2.

Accordingly, $R$ is represented by
\begin{align*}
R &= \frac{T_{2k}(x)\Lambda_{(4k+1,1)}(x+(2k+1)i)}{\phi(x+(6k+\frac{9}{2})i)\phi(x-(2k+\frac{3}{2})i)\phi(x-(6k+\frac{1}{2})i)} \\
&\quad \times \left( T_1(x-(10k+5)i)\Lambda_{(2k,2k)/(2k-1)}(x-(8k+2)i) \\
&\quad - T_0(x-(10k+6)i)\Lambda_{(2k-1,2k-1)/(2k-2)}(x-(8k+1)i) \right).
\end{align*}

The second term in the rhs coincides with that in the bracket of (E.16). Thus
\begin{align*}
R &= \frac{T_{2k}(x)t_{1}^{(6)}(x+\frac{3i}{2k+4})}{\phi(x-(6k+\frac{1}{2})i)\phi(x-(10k+\frac{15}{2})i)} \Lambda_{(2k+1,2k+1,1)/(2k)}(x-(8k+4)i)
\end{align*}

where $\phi(x-(10k+\frac{7}{2})i) = \phi(x+(6k+\frac{9}{2})i)$ and the definition of $t_{1}^{(6)}$ is used. Finally, we apply lemma 3 to the above and reach expression (E.15). Thereby we prove the functional relation.
Appendix F. A result for an exceptional case

We sketch the result for exceptional cases, $L = 3 \sim 6$.

For $L = 3, 4, 6$, the $Y$-system is neatly described by

$$Y_j \left( x \pm \frac{L - 2}{2L + 4} \right) = \prod_{j-t} \Xi \epsilon(x) \quad (j = 1, \ldots, r)$$

where the corresponding Dynkin diagrams are those for $E_r$, and $r = 8, 7, 6$, respectively.

The solutions of $Y$ in terms of transfer matrices have been reported in [22] for $L = 3$ and [23] for $L = 4, 6$.

The case $L = 5$, corresponding to $M_{5,6} + \phi_{1,2}$, is a little complicated as announced in [73]. Below, we shall summarize the result then supplement the outline of the proof as promised there.

In this case, there are more breathers in the system than the cases we discussed above. Therefore, we need to introduce more $s l_2$ origin objects.

$$T_{B_1}(x) := T_1(x) \quad (F.1)$$

$$T_{B_3}(x) := \Lambda_{(8,1)}(x + \frac{12}{14} i)/\phi(x - \frac{12}{14} i) \quad (F.2)$$

$$T_{B_5}(x) := \Lambda_{(15,8,8)/(7,7)}(x)/\phi(x + \frac{3}{2} i) \quad (F.3)$$

$$T_{B_7}(x) := \Lambda_{(15,15,8,8)/(14,7,7)} \left( x + \frac{12}{14} i \right)/\phi(x - \frac{12}{14} i) \phi(x + \frac{3}{2} i)) \quad (F.4)$$

$$T_{B_2}(x) := T_7(x) \quad (F.5)$$

$$T^{(6)}(x) := \Lambda_{(8,7)/(6)}(x + \frac{25}{14} i); \quad (F.6)$$

then the following relations hold.

$$T_{B_1}(x + \frac{3}{14} i) = T_0(x + \frac{14}{14} i) + \phi(x - \frac{12}{14} i)T_{B_3}(x)$$

$$T_{B_3}(x + \frac{3}{14} i) = T_0(x)T_0(x + \frac{8}{14} i) + T_{B_1}(x)T_{B_5}(x)$$

$$T_{B_7}(x + \frac{3}{14} i) = T_0(x)T_0(x + \frac{14}{14} i) + T_{B_3}(x)T_{B_7}(x)$$

$$T_{B_2}(x + \frac{3}{14} i) = T_0(x + \frac{14}{14} i) + T^{(6)}(x).$$

For a magnon-like $t$-system we assume

$$t^{(1)}_1 \left( x + \frac{i}{14} \right) t^{(1)}_1 \left( x - \frac{i}{14} \right) = t^{(1)}_2(x) + t^{(3)}_1(x)g^{(1)}_1(x) \quad (F.7)$$

$$t^{(1)}_2 \left( x + \frac{i}{14} \right) t^{(1)}_2 \left( x - \frac{i}{14} \right) = t^{(1)}_1(x)t^{(1)}_3(x) + t^{(3)}_2(x)g^{(1)}_2(x) \quad (F.8)$$

$$t^{(1)}_3 \left( x + \frac{i}{14} \right) t^{(1)}_3 \left( x - \frac{i}{14} \right) = t^{(1)}_2(x)t^{(1)}_4(x) \quad (F.9)$$

$$t^{(2)}_1 \left( x + \frac{i}{14} \right) t^{(2)}_1 \left( x - \frac{i}{14} \right) = t^{(2)}_2(x) + t^{(3)}_1(x)g^{(2)}_1(x) \quad (F.10)$$

$$t^{(2)}_2 \left( x + \frac{i}{14} \right) t^{(2)}_2 \left( x - \frac{i}{14} \right) = t^{(2)}_1(x)t^{(2)}_3(x) + t^{(3)}_2(x)g^{(2)}_2(x) \quad (F.11)$$

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Remark. They are obtained from the $T$-system for $D_4$, by requiring $t_3^{(3)} = t_3^{(4)} = 0$.

The identification between $Y$ and $T$ is given by

\begin{align}
Y_{B_1} &= \phi(x - \frac{12}{7}i) T_{B_3}(x)/T_0(x \pm \frac{11}{14}i) \\
Y_{B_3} &= T_{B_1}(x) T_{B_5}(x)/(T_0(x) T_0(x \pm \frac{3i}{7})) \\
Y_{B_5} &= T_{B_3}(x) T_{B_7}/(T_0(x \pm \frac{3}{14}i) T_0(x \pm \frac{5}{14}i)) \\
Y_{B_2} &= T^{(6)}(x)/T_0(x \pm \frac{1}{14}i) \\
Y_{K_1}(x) &= \frac{t_2^{(1)}(x)}{t_3^{(3)}(x) g_1^{(1)}(x)} \\
Y_{K_2}(x) &= \frac{t_2^{(2)}(x)}{t_3^{(3)}(x) g_1^{(2)}(x)} \\
Y^{(1)}(x) &= \frac{t_2^{(3)}(x)}{t_1^{(3)}(x) t_2^{(2)}(x) t_1^{(4)}(x)} \\
Y^{(2)}(x) &= \frac{t_2^{(4)}(x)}{t_1^{(3)}(x) t_1^{(4)}(x)}.
\end{align}

The explicit solution is asserted as follows. The following choice of $t_{1,2}^{(a)}(x)$ solves the kink $T$-system:

\begin{align}
t_1^{(1)}(x) &= T^{(6)}(x) \\
t_1^{(2)}(x) &= T_{B_7}(x) \\
t_1^{(3)}(x) &= \Lambda_{(12,8,7)/(5,4)}(x) \\
t_1^{(4)}(x) &= \Lambda_{(5,1)}(x + \frac{15}{14}i)/\phi(x - \frac{13}{2}i + \frac{15}{14}i) \\
t_2^{(1)}(x) &= T_{B_5}(x) T_{B_2}\left(x \pm \frac{i}{7}\right)
\end{align}
Lemma 15 $T_{B_5}(x) = T_{B_5}(x \pm \frac{i}{7})T_{B_2}(x)$ (F.22)
Lemma 14 $T_{B_3}(x) = t^{(3)}_2(x)$ with additional constraints
$T_{B_3} = \frac{t^{(2)}_3(x)}{t^{(3)}_2(x)}$
$g^{(1)}_1(x) = T_0(x), \quad g^{(2)}_1(x) = T_0(x \pm \frac{2i}{7})$

By these choices the above $Y$ satisfy the $Y$-system proposed in [25],
$Y_{B_1}(x \pm \frac{3}{14}i) = \Xi_{B_3}(x)$
$Y_{B_3}(x \pm \frac{3}{14}i) = \Xi_{B_1}(x)\Xi_{B_5}(x)$
$Y_{B_5}(x \pm \frac{3}{14}i) = \Xi_{B_3}(x)\Xi_{K_2}(x \pm \frac{2}{14}i)\Xi_{K_1}(x)\Xi^{(1)}(x \pm \frac{1}{14}i)\Xi^{(2)}(x)$
$Y_{B_2}(x \pm \frac{3}{14}i) = \Xi_{K_1}(x \pm 2)\Xi^{(1)}(x \pm \frac{1}{14}i)\Xi_{K_2}(x)\Xi^{(2)}(x)$
$Y_{K_3}(x \pm \frac{1}{14}i) = \Xi_{B_5}(x)\mathcal{L}^{(1)}(x)$
$Y_{K_1}(x \pm \frac{1}{14}i) = \Xi_{B_3}(x)\mathcal{L}^{(1)}(x)$
$Y^{(1)}(x \pm \frac{1}{14}i) = \mathcal{L}^{(2)}(x)\mathcal{L}_{K_2}(x)\mathcal{L}_{K_1}(x)$
$Y^{(2)}(x \pm \frac{1}{14}i) = \mathcal{L}^{(1)}(x)$.

We first give a few lemmas concerning equivalent expressions for transfer matrices.

**Lemma 14** $T_{B_5}(x)$ has the following two equivalent expressions in addition to (F.3).
$T_{B_5}(x) = T_{D_2}(x) = \frac{1}{\phi(x \pm \frac{3}{14}i)}\Lambda_{(8,8,1)/(7)}(x + \frac{12}{7}i)$.

The first equality has already appeared in (25). The second is shown by the duality.

**Lemma 15** $T_{B_7}(x) = t^{(2)}_1(x)$ is also written as follows.
$T_{B_7}(x) = \frac{1}{\phi(x + \frac{22}{14}i)}\Lambda_{(4,3,1)/(2)}(x - \frac{31}{14}i)$.

To prove this, we utilize $T_{(15,15,8,8)/(14,7,7)}(x)$. The quantum Jacobi–Trudi formula and the duality relation conclude
$T_{(15,15,8,8)/(14,7,7)}(x) = \Lambda_{(15,15,8,8)/(14,7,7)}(x)$
$+ T_0(x)T_0(x - 2i)(T_1(x + 17i)T_4(x - i) - T_0(x + 16i)T_3(x))$. (F.25)
On the other hand, the same quantity is given by

$$T_{(15,15,8,8)/(14,7,7)}(x) = \phi \left( x + \frac{i}{2} \right) \phi \left( x + \frac{13i}{14} \right) T_7(x - 11i) \Lambda_{(8,8)/(7)}(x + 10i)$$  \hfill (F.26)

thanks to the semi-standard condition. Again, by the quantum Jacobi–Trudi formula and the duality relation, $\Lambda_{(8,8)/(7)}(x + 10i) = T_1(x - 7i)T_2(x + \frac{3i}{7}) - T_0(x - \frac{3i}{7})T_1(x - 2i)$. Using this expression, as well as with some rearrangement, we find

$$T_7(x - 11i) \Lambda_{(8,8)/(7)}(x + 10i) = T_1(x - 7i)T_{D_2}(x - i) - T_0(x - \frac{3i}{7}) \Lambda_{(2,1)}(x)$$

$$+ T_0(x - i)(T_4(x - i)T_1(x - 7i) - T_0(x - \frac{3i}{7})T_3(x)).$$

Substitute this into (F.26), compare the resultant expression with (F.25) and we have

$$\Lambda_{(15,15,8,8)/(14,7,7)}(x) = \phi \left( x + \frac{i}{2} \right) \phi \left( x + \frac{13i}{14} \right)$$

$$\times \left( T_1(x - 7i)T_{D_2}(x - i) - T_0 \left( x - \frac{3i}{7} \right) \Lambda_{(2,1)}(x) \right).$$

By the representation of $T_1(x - 7i)$ by a $2 \times 1$ table, the diagrammatical rule immediately leads to

$$\Lambda_{(15,15,8,8)/(14,7,7)}(x) = \frac{d_{B_2}(x - \frac{11i}{2})}{\phi(x + \frac{11i}{2})} \Lambda_{(4,3,1)/(2)}(x - 3i),$$

where $d_{B_2}(x)$ specifies the denominator in (F.28). This proves the lemma after a shift in $x$ and the simplification in the coefficient.

**Lemma 16** $t_1^{(3)}(x)$ can be also represented by

$$t_1^{(3)}(x) = \frac{1}{\phi(x \pm \frac{3i}{2})} \Lambda_{(7,7,5,1)/(6,2)} \left( x + \frac{12}{7}i \right).$$

Just as in the proof of lemma 15, two expressions for $T_{(12,8,7)/(5,4)}(x)$ lead to

$$t_1^{(3)}(x) = \Lambda_{(12,8,7)/(5,4)}(x) = \frac{1}{\phi(x \pm \frac{3i}{2})} (\Lambda_{(5,1)}(x + 8i) \Lambda_{(5,5)/(4)}(x + 16i)$$

$$- \phi(x \pm \frac{3i}{2}) T_7(x) T_{D_2}(x)).$$  \hfill (F.27)

The product $\Lambda_{(5,1)}(x + 8i) \Lambda_{(5,5)/(4)}(x + 16i)$ decomposes as follows.

$$(\Lambda_{(5,1)}(x + 8i) \Lambda_{(5,5)/(4)}(x + 16i)) = (\Lambda_{(5,1)}(x - 4i) \Lambda_{(5,5)/(4)}(x + 4i))_{x->x+(12/7)i}$$

$$= \Lambda_{(7,7,5,1)/(6,2)}(x + \frac{12}{7}i) + T_2(x + \frac{12}{7}i) \Lambda_{(8,8,1)/(7)}(x + \frac{12}{7}i)$$

$$= \Lambda_{(7,7,5,1)/(6,2)}(x + \frac{12}{7}i) + T_7(x) T_{B_2}(x) \phi(x \pm \frac{3i}{2})$$  \hfill (F.28)

where we have used lemma 14 in the last equality. Using (F.28) in (F.27) and remembering $T_{B_2}(x) = T_{D_2}(x)$, one easily establishes the lemma.

**Lemma 17** The following equality can easily be verified:

$$\Lambda_{(6,6,4,3,1)/(5,3,2)}(x) = \phi \left( x \pm \frac{3}{14}i \right) T_{B_2} \left( x \pm \frac{i}{7} + \frac{12}{7}i \right).$$

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The diagrammatic decomposition yields
\[ \Lambda_{(6,6,4,3,1)/(5,3,2)}(x) = \frac{1}{f_2(x)}\Lambda_{(3,3,1)/(2)}(x \pm 5i). \]

The \( m = 3 \) case of lemma 3 and \( T_{B_5}(x) = T_{D_2}(x) \) then validates the lemma.

With these preparations, we prove the magnon-like \( t \)-system.

**Proof of (F.10).** With (F.17), (F.6) and using the duality and the quantum Jacobi–Trudi formula, we can express the product \( t_1^{(1)}(x \pm \frac{1}{14}i) \) using only \( T_0(x) \) and \( T_2(x) \). Similarly, \( t_1^{(3)} \) can be rewritten by \( T_0(x), T_2(x) \) and \( T_4(x) \), according to (F.19). Then it is immediately seen that
\[ T_0(x)t_1^{(3)}(x) - t_1^{(1)}(x \pm \frac{1}{14}i) = T_2(x \pm \frac{11}{7}i)T_{D_2}(x). \]

Then the equality \( T_{B_5}(x) = T_{D_2}(x), T_2(x + \frac{12}{7}i) = T_7(x) \) and the definition of \( t_2^{(1)}(x) \) in (F.21) concludes (F.7).

**Proof of (F.11).** Consider the product \( \Lambda_{(4,3,1)/(2)}(x - 4i)\Lambda_{(4,4,2)/(3,1)}(x + 4i) \) which is equal to
\[ \phi(x \pm \frac{3}{14}i)t_1^{(2)}(x \pm \frac{1}{14}i + \frac{12}{7}i) \] (F.29)
due to lemma 15.

As usual, the diagrammatical argument leads to
\[ \Lambda_{(4,3,1)/(2)}(x - 4i)\Lambda_{(4,4,2)/(3,1)}(x + 4i) = \phi(x \pm \frac{3}{2}i)\phi(x \pm \frac{5}{2}i)\Lambda_{(7,7,5,1)/(6,2)}(x) + T_2(x)\Lambda_{(6,6,4,3,1)/(5,3,2)}(x). \]

Then we apply lemmas 16 and 17 to the first and the second term in the rhs, respectively, to reach
\[ \phi(x \pm \frac{3}{14}i)(t_1^{(3)}(x + \frac{12}{7}i)T_0(x \pm \frac{2}{7}i + \frac{12}{7}i) + T_{B_5}(x + \frac{12}{7}i)T_{B_5}(x \pm \frac{1}{7}i + \frac{12}{7}i). \] (F.30)

By equating (F.29) and (F.30), we show the validity of (F.8).

**Proof of (F.13).** This is actually equivalent to (F.27). Note that \( \Lambda_{(5,1)}(x + 8i) = \phi(x + \frac{3}{2}i)t_1^{(4)}(x + \frac{1}{14}i) \) and that the diagrammatic symmetry leads to \( \Lambda_{(5,5)/(4)}(x - 8i) = \phi(x - \frac{3}{2}i)t_1^{(4)}(x - \frac{1}{14}i) \). Note also \( T_7(x) = T_{B_2}(x) \). Then one easily verifies (F.10) using \( t_2^{(4)}(x) \) in (F.24).

**Proof of (F.16).** This is the final non-trivial relation. This relation can be simply shown by using the result for the \( A_{4k+1} \) case. We simply put \( k = 1 \) and identify \( t_1^{(1)}, t_1^{(2)}, t_1^{(3)}, t_1^{(4)}, t_2^{(3)}, g_1^{(1)}, g_1^{(2)} \) here with \( t_1^{(3)}, t_1^{(4)}, t_1^{(5)}, t_1^{(6)}, t_2^{(5)}, t_1^{(1)}, t_1^{(2)} \) there and use the duality \( T_2 = T_7^* \).

All remaining relations can be checked easily. Thus the above argument completes the proof of the magnon-like \( t \)-system of the dilute \( A_5 \) model.
Appendix G: Numerical results for the dilute $A_8$ model in regime 2

We supplement some numerical data supporting the conjecture of the analyticity of $Y$-functions in section 11.1. Throughout this appendix, we choose $\beta = 0.4$ and $J = 1$.

The solution to the Bethe ansatz equation, corresponding to the largest eigenvalue of QTM, assumes the form of two strings. We list the explicit locations of roots for $N = 10$ and $q = 0$.

| BAE Roots $q = 0$ |
|-------------------|
| $0.13548013966460354 \pm 0.4464809553245548i$ |
| $0.03263577725129405 \pm 0.4935687693509952i$ |
| $0.002550182914375612 \pm 0.4961489632802383i$ |
| $-0.02495383965730369 \pm 0.4946627475810687i$ |
| $-0.09400161546158932 \pm 0.47389905045124336i$ |

Their imaginary parts are approximately equal to $\pm \frac{1}{2}$, irrespective of $L, N, q$ and $\beta$ (see figure G.1).

**Figure G.1.** The Bethe ansatz roots for $N = 10$ for $q = 0, 0.1, 0.2, 0.3, 0.4$ and 0.5. Their location changes only slightly.

**Figure G.2.** Zeros of $T_1(x)$ (left) and $T_2(x)$ (right) for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4$ and 0.5.
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

**Figure G.3.** Zeros of $T_3(x)$ (left) and $T_4(x)$ (right) for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4$ and $0.5$.

**Figure G.4.** Zeros of $T_5(x)$ (left) and $T_6(x)$ (right) for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4$ and $0.5$.

**Figure G.5.** Zeros of $T_7(x)$ for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4$ and $0.5$.

The zeros for the symmetric fusion QTM, $T_m$, ($m = 1, \ldots, 7$), are plotted in figures G.2–G.5. These figures clearly show that the simple-minded choice of $Y$ in (17) is not appropriate: the choice requires that $T_m$ should be analytic in the strip $\text{Im } x \in [-1, 1]$. 

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The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

Figure G.6. Zeros of $T_{B_3}(x)$ (left) and $T_{B_5}(x) (= T_{D_2})$ (right) for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4, 0.5$.

Figure G.7. Zeros of $t_1^{(0)}(x)$ (left) and $t_1^{(2)}(x)$ (right) for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4$ and 0.5.

(Or the arguments can be simultaneously shifted by a half period so that the strip is $\text{Im} \ x \in [0.8, 1.8] \cup [-1.8, -0.8]$. ) Obviously, neither choice of strip results in a set of $T_m$ free from zeros. (In particular, $T_3$ and $T_5$ break the analyticity in both strips.) Therefore, the transformation of the $Y$-system to the TBA is not legitimate.

The difference in the argument $x$ appearing in the breather $T$, $s l_3 \ T$ and magnon-like $t$ is much smaller than unity, which is crucial.

$T_{B_1}(x)$ is identified with $T_1(x)$. The functional relation then implies that it should be analytic within $\text{Im} \ x \in \left[-\frac{4}{10}, \frac{4}{10}\right]$. This can easily be checked from figure G.2. Among other $T_m$, $T_4(x)$ and $T_2(x + \frac{6}{5}i)$ are identified with $t_1^{(1)}(x)$ and $t_2^{(3)}(x)$. Their ANZC property in a wider strip than $\text{Im} \ x \in \left[-\frac{1}{5}, \frac{1}{5}\right]$ is clearly seen.

The breather related QTM, $T_{B_3}(x)$, is equal to $t_1^{(7)}(x)$, and $T_{B_5}(x) (= T_{D_2}(x))$ is $t_2^{(4)}(x)$, thus the information on their zeros is significant. In particular, the latter appears frequently in the solution of other $t_m^{(a)}$ ($m = 2, 3$).

Their zeros are shown in figure G.6.

The remaining zeros of QTM to be specified are those for $t_1^{(0)}(x), t_1^{(2)}(x), t_1^{(3)}(x)$ and $t_1^{(4)}(x)$. They are depicted in figures G.7 and G.8, respectively.
The dilute $A_L$ models and the $\Phi_{1,2}$ perturbation of unitary minimal CFTs

By combining the above data on zeros, one can verify that all $Y$ and $1 + Y$ possess no zeros or singularities from QTM in their own strips. Only trivial zeros (or singularity) of order $N$ exist for $Y_{B_i}$ simply coming from the normalization factor $\phi^r(x)$ (or $T_0(x)$).

Appendix H. The kernel matrices for $L = 5, 7$

For $L = 5$, we define the order as

$$t(\hat{\log} Y_{B_1}, \hat{\log} Y_{B_5}, \hat{\log} Y_{B_5}, \hat{\log} Y_{B_2}, \hat{\log} Y_{K_1}, \hat{\log} Y^{(1)}, \hat{\log} Y^{(2)})$$

then

$$\hat{M} = \begin{pmatrix}
2c h \frac{3}{14} k & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\
-1 & 2c h \frac{3}{14} k & -1 & 0 & 0 & 0 & 0 & 0 \\
0 & -1 & 2c h \frac{3}{14} k & 0 & -1 & 2c h \frac{3}{14} k & 2c h \frac{1}{14} k & -1 \\
0 & 0 & 0 & 2c h \frac{3}{14} k & -2c h \frac{2}{14} k & -1 & -2c h \frac{1}{14} k & , -1 \\
0 & 0 & 0 & -1 & 2c h \frac{1}{14} k & 0 & 0 & 0 \\
0 & 0 & -1 & 0 & 0 & 2c h \frac{1}{14} k & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 2c h \frac{1}{14} k & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 2c h \frac{1}{14} k & 0
\end{pmatrix}$$

$$\hat{K}_0 = \begin{pmatrix}
0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 1 & 2c h \frac{2}{14} k & 2c h \frac{1}{14} k & , 1 \\
0 & 0 & 0 & 0 & 2c h \frac{2}{14} k & 1 & 2c h \frac{1}{14} k & , 1 \\
0 & 0 & 0 & 1 & 0 & 0 & -1 & 0 \\
0 & 0 & 1 & 0 & 0 & 0 & -10 & 0 \\
0 & 0 & 0 & -1 & -1 & 0 & -1 & -1 \\
0 & 0 & 0 & 0 & 0 & 0 & -1 & 0
\end{pmatrix}$$

Figure G.8. Zeros of $t_1^{(3)}(x)$ (left) and $t_1^{(4)}(x)$ (right) for $N = 10$ and $q = 0.1, 0.2, 0.3, 0.4$ and $0.5$.  

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For $L = 7$, we define the order as
\[
\left( \log Y_{B_1}, \log Y_{B_3}, \log Y_{0}^{(1)}, \log Y_{0}^{(2)}, \log Y_{1}^{(1)}, \log Y_{1}^{(2)}, \log Y_{1}^{(3)}, \log Y_{1}^{(4)}, \log Y_{1}^{(5)}, \log Y_{1}^{(6)} \right)
\]
then
\[
\hat{M} = \begin{pmatrix}
2 \operatorname{ch} \frac{1}{18} k & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
-1 & 2 \operatorname{ch} \frac{1}{18} k & 0 & -2 \operatorname{ch} \frac{1}{18} k & 0 & -1 & 0 & 0 & 0 \\
0 & 0 & 2 \operatorname{ch} \frac{1}{18} k & -1 & 0 & 0 & 0 & 0 & 0 \\
0 & -1 & 2 \operatorname{ch} \frac{1}{18} k & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 2 \operatorname{ch} \frac{1}{18} k & 0 & -1 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 2 \operatorname{ch} \frac{1}{18} k & 0 & -1 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 2 \operatorname{ch} \frac{1}{18} k & -1 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 \operatorname{ch} \frac{1}{18} k & -1 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 \operatorname{ch} \frac{1}{18} k
\end{pmatrix}
\]
\[
\hat{K}_0 = \begin{pmatrix}
0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
1 & 0 & 0 & 2 \operatorname{ch} \frac{1}{18} k & 0 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & -2 \operatorname{ch} \frac{1}{18} k & 0 & -2 \operatorname{ch} \frac{1}{18} k & -1 & -2 \operatorname{ch} \frac{1}{18} k & -1 \\
0 & 1 & 1 & 0 & -2 \operatorname{ch} \frac{1}{18} k & -1 & -2 \operatorname{ch} \frac{1}{18} k & -2 \operatorname{ch} \frac{1}{18} k & -1 \\
0 & 0 & -1 & 0 & 0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & -1 & 0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 1 \\
0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0
\end{pmatrix}.
\]

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