Observation of Exclusive Deeply Virtual Compton Scattering in Polarized Electron Beam Asymmetry Measurements

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We report the first results of the beam spin asymmetry measured in the reaction $\vec{e}p \rightarrow e p\gamma$ at a beam energy of 4.25 GeV. A large asymmetry with $a \sin\phi$ modulation is observed, as predicted for the interference term of Deeply Virtual Compton Scattering and the Bethe-Heitler process. The amplitude of this modulation is $a = 0.202 \pm 0.028$. In leading order and leading twist pQCD, the $a$ is directly proportional to the imaginary part of the DVCS amplitude.

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High-energy lepton nucleon scattering experiments revealed the quark structure of the nucleon and were employed to measure the longitudinal quark momentum distributions as well as the quark spin distributions. The recently developed formalism \[3\] for the QCD description of deeply exclusive leptoproduction reactions introduces Generalized Parton Distributions (GPDs). They carry new information about the dynamical degrees of freedom inside the nucleon not accessible in inclusive experiments. It has been shown in the Bjorken scaling regime (high photon virtuality $Q^2$, high energy transfer at fixed Bjorken scaling variable $x_B$) the scattering amplitude for exclusive processes can be factorized into a hard-scattering part (exactly calculable in perturbative QCD) and a nucleon structure part parametrized via GPDs. This is called the “handbag approximation”, depicted in Fig. 1a. In addition to the dependence on the parton momentum fraction $x$ GPDs depend on two more parameters, the skewedness $\xi$, and the momentum transfer $t$ to the baryonic system. In the Bjorken regime $\xi \rightarrow x_B/(2-x_B)$, and is related to the momentum imbalance between the struck quark and the quark that is put back into the final state baryon. The momentum transfer $t$, which can be varied independently in exclusive processes, measures the transverse momentum transfer between interacting quarks. The GPDs reduce to the longitudinal parton distributions in the limit $\xi \rightarrow 0$, $t \rightarrow 0$.

Deeply exclusive experiments probe the nucleon wave function at the quark-gluon level, expressed in terms of two spin-dependent and two spin-independent GPDs. Deeply Virtual Compton Scattering (DVCS), i.e. electroproduction of photons from nucleons in the Bjorken regime is most suitable for studying GPDs at moderately high energies. While exclusive meson production requires high energies and high photon virtualities to reach the Bjorken regime, DVCS may access the GPDs already at $Q^2$ as low as 1 (GeV/c)$^2$ \[4\].

In this letter we present the first observation of the fully exclusive DVCS signal in the beam spin asymmetry measured in the reaction $\vec{e}p \rightarrow e p\gamma$. The measurement was performed using a 4.25 GeV longitudinally polarized electron beam and the CEBAF Large Acceptance Spectrometer (CLAS) \[7\] at the Thomas Jefferson National Accelerator Facility. Observation of the beam-spin asymmetry in semi-exclusive production of high energy photons has been reported recently by the HERMES collaboration \[5\], and cross section measurements of DVCS by the H1 collaboration at DESY at very small $x_B$ \[9\], which is complementary to the valence quark regime of this measurement.

The experiments measures the DVCS contribution through its interference with the Bethe-Heitler process. In contrast to DVCS, where the photon is emitted from the nucleon, BH photons are emitted from the incoming or scattered electrons (Fig.1). While the BH cross section in most of the kinematic region is much larger than the DVCS cross section, the interference between the two amplitudes boosts the effect of DVCS and produces a large cross section difference for electron helicities aligned or anti-aligned with the electron beam. In this difference, the large helicity-independent BH contribution drops out, and only the helicity-dependent interference term remains \[10\]:

$$
\frac{d^4\sigma^+}{dQ^2 dx_B dt d\phi} - \frac{d^4\sigma^-}{dQ^2 dx_B dt d\phi} \\
\propto Im (T^{DVCS}) T^{BH} \propto a \cdot Im\hat{M}^{1,1} \sin\phi \\
+ b \cdot Im\hat{M}^{0,1} \sin 2\phi + O\left(\frac{1}{Q^2}\right) + \cdots .
$$

FIG. 1. Feynman diagrams for DVCS and Bethe-Heitler processes contributing to the amplitude of $ep \rightarrow ep\gamma$ scattering.

Here $a$ and $b$ are independent of $\phi$, the angle between the lepton and hadron planes, and $\hat{M}^{1,1}$ and $\hat{M}^{0,1}$ are the helicity amplitudes for transverse and longitudinal virtual photons, respectively. The cross section for posi-
tive (negative) electron beam helicity is denoted with \( \sigma^+ (\sigma^-) \). This difference measures directly the imaginary part of the DVCS amplitude, since the BH amplitude can be calculated exactly. In leading-order only \( M^{1,1} \) contributes [4].

Experimentally it is much simpler to measure an asymmetry \( A = (d^3\sigma^+ - d^3\sigma^-)/(d^3\sigma^+ + d^3\sigma^-) \) rather than a cross section difference. This asymmetry gives the relative strength of the relevant amplitudes and has a more complex dependence on \( \phi \) than the cross section difference, due to \( \phi \) terms in the denominator. In the Bjorken limit, however, the \( \phi \) dependence of the denominator is expected to be small.

In this analysis, we used CLAS electroproduction data taken in March 1999. Scattering of 4.25 GeV longitudinally polarized electrons from a liquid hydrogen target was studied in a wide range of kinematics. The analyzed event sample corresponds to an integrated luminosity of 1.3 fb\(^{-1}\).

The reaction \( ep \rightarrow ep\gamma \) is identified by analyzing the missing mass squared distributions in the reaction \( ep \rightarrow epX \): \( M_x^2 = (\nu + M - E_p)^2 - (\vec{q} - \vec{P}_p)^2 \), where \( \nu \) and \( \vec{q} \) are the virtual photon energy and momentum, \( E_p \) and \( \vec{P}_p \) are the energy and the momentum of the recoil proton, and \( M_x \) is the mass of the proton.

The main background to the single photon final state is from \( \pi^0 \) production. At large missing energy the missing mass resolution of CLAS is not sufficient to separate single photons and \( \pi^0 \)s event-by-event. The number of single photon events was determined using a fitting technique that analyzes the line shape of the \( M_x^2 \) distributions. As a fit function the sum of two Gaussians and a third polynomial is used:

\[
F(M_x^2) = N_\gamma \cdot e^{-(M_x^2 - M_{\gamma}^2)/2\sigma_\gamma^2} + N_{\pi^0} \cdot e^{-(M_x^2 - M_{\pi^0}^2)/2\sigma_{\pi^0}^2} + \sum_{j=0}^{3} P_j \cdot (M_x^2)^j. \tag{2}
\]

The Gaussians represent the \( M_x^2 \) distributions for single photon and single pion final states. The polynomial is used to model the smooth background that arises from radiative processes. The parameters \( M_{\gamma}^2 \) and \( \sigma_{\gamma}^2 \) are determined from the fit to the \( M_x^2 \) distribution of selected Bethe-Heitler events (Figure 2.a). For \( M_{\pi^0}^2 \) and \( \sigma_{\pi^0}^2 \), the results of a similar fit to the \( M_x^2 \) distribution of selected events with final state \( ep \rightarrow ep\gamma\pi^0 \) are used (Figure 2.b). In these events both photons are detected in the CLAS electromagnetic calorimeter; they have an invariant mass near the \( \pi^0 \) mass. In Figure 2 the curves represent the fit to a sum of the Gaussian and a third order polynomial. The samples of BH and \( \pi^0 \) events were selected in the same kinematical range as the main sample for DVCS studies.

The \( M_x^2 \) distributions are analyzed in eleven \( \phi \) bins. The final fit to the \( M_x^2 \) distributions with the function in Eq. 2 is performed separately for each helicity state and for the sum. In the fit only two parameters are adjusted: the number of single photon \( (N_\gamma) \) and the number of \( \pi^0 \) \( (N_{\pi^0}) \) events. For each \( \phi \) bin the shape of the background was fixed from a fit to the points in the summed \( M_x^2 \) distribution that excluded the \( M_x^2 \) range from -0.07 GeV\(^2\) to 0.08 GeV\(^2\). The relative magnitude of the background for distributions at positive and negative helicities are found from the fit to the same \( M_x^2 \) range as for the sum.

Figure 3 shows a typical fit result for the sum of two helicities at \( \phi = 90^\circ \). The solid line corresponds to the fitted function \( F(M_x^2) \), the dashed line represents the Gaussian for the missing photon, and the dashed-dotted line is the Gaussian for the missing \( \pi^0 \). The dotted line represents the polynomial background. In this bin the number of photon and pion events are \( N_\gamma = 4201 \pm 109 \) and \( N_{\pi^0} = 2010 \pm 101 \), respectively. The \( \chi^2 \) per degree of freedom of the fit is 1.8. Similar results have been obtained for all kinematical bins.

The fitted number of single photon events are used to calculate the beam spin asymmetry as:

\[
A = \frac{1}{P_e} \frac{(N_+ - N_-)}{(N_+ + N_-)}. \tag{3}
\]
where $P_e$ is the beam polarization, $N_{\gamma}^{\pm(-)}$ is the number of $ep \rightarrow ep\gamma$ events at positive (negative) beam helicity. The average beam polarization, $P_e = 70\%$, was measured using Möller scattering.

In Figure 3 we present our main result, the $\phi$ dependence of $A$. Data in each $\phi$ bin are integrated in the range of $Q^2$ from 1.00 (GeV/c)$^2$ to 1.75 (GeV/c)$^2$ and $-t$ from 0.1 (GeV/c)$^2$ to 0.3 (GeV/c)$^2$. The error bars shown are statistical. Most of the systematic uncertainties related to the experiment do not contribute to $A$. Only the error in the measurement of beam polarization, $\pm 1.65\%$, remains. There is also a systematic error in the calculation of $N_{\gamma}$ due to the determination of the mean and standard deviation of the Gaussian functions for the photon and pion missing mass squared distributions, and also due to the fit procedure to the $M^2_x$ distributions. These errors are defined as a deviation of $A$ from its central value, when the mean and standard deviation of Gaussians are shifted within their errors, and when different fit techniques were used (see [12] for details).

The data points are fitted with the function $A(\phi) = \alpha \sin \phi + \beta \sin 2\phi$. The fitted parameters are $\alpha = 0.202 \pm 0.028_{\text{stat}} \pm 0.013_{\text{sys}}$ and $\beta = -0.024 \pm 0.021_{\text{stat}} \pm 0.009_{\text{sys}}$. In the Bjorken regime $\beta$ should vanish, leaving only the contribution from transverse photons (see e.g. Ref. [10]). In Figure 3 the dark shaded region corresponds to the range of the fitted function within the statistical uncertainties of $\alpha$ and $\beta$. The light shaded region includes systematic uncertainties on these parameters, estimated using the method described above.

The resulting asymmetry is in a good agreement with a $\sin \phi$ modulation. Curves in Figure 4 show the results of theoretical calculations from Refs. [6,11,13,14] at fixed values of $Q^2 = 1.25$ (GeV/c)$^2$, $x_B = 0.19$, and $-t = 0.19$ (GeV/c)$^2$. The limited experimental information does not allow to unambiguously extract GPDs from the measurement. Description of GPDs that model the $\xi$ and $t$ dependencies are therefore used to predict observables accessible in experiments. The dashed curve is a calculation at leading-twist [5] and no $\xi$ dependence in the evaluation of GPDs, the dotted curve is leading-twist with $\xi$ dependence [5], and the solid curve includes twist-3 effects. All three calculations include the $D$-term in the parameterization of the GPDs [10], which is related to double pion contributions. For a more detailed description of the model assumptions we refer to a recent review [11]. We have estimated that the model asymmetries would be reduced by about 7% if they are averaged over the experimental acceptances, bringing them somewhat closer to the measured data points.

Although the experimental results are close to the lower range of theoretical predictions, none of the calculations is in agreement with our data. This could be due to several factors which need to be studied in future research. Firstly, parameterizations of the $\xi$ dependence of the GPDs were modeled in only a few different
ways. Secondly, higher order perturbative QCD corrections have not been included. They tend to reduce the measured asymmetry \[17\]. Thirdly, it is possible that the average \(Q^2\) and \(W\) of our measurements are not sufficiently high to fully reach the Bjorken regime, so that some non-perturbative corrections may be needed. This can be tested in a future experiment by measuring DVCS at the same \(x_B\) and \(t\) but at higher \(Q^2\) and \(W\). Such experiments have been approved and are currently in preparation \[18,19\].

In conclusion, we have presented the first measurement of the beam spin asymmetry in exclusive electroproduction of real photons in the deeply inelastic regime. We see a clear asymmetry, as expected from the interference of the DVCS and BH processes (Eq.1). Our results agree in sign, and are not far in magnitude to predictions based on available models of GPD parameterizations. This strongly supports expectations that DVCS will allow access to GPDs at relatively low energies and momentum transfers. This opens up a new avenue for the study of nucleon structure, which is inaccessible in inclusive scattering experiments. Further measurements at higher beam energy are planned, which will allow significant expansion of the \(Q^2\), \(x_B\), and \(t\) dependences of GPDs.

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