Is there a connection between “dark” and “light” physics?

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Abstract. In the early-mid 20th century Dirac and Zel’dovich were among the first scientists to suggest an intimate connection between cosmology and atomic physics. Though a revolutionary proposal for its time, Dirac’s Large Number Hypothesis (1937) adopted a standard assumption of the day, namely, the non-existence of the cosmological constant term ($\Lambda = 0$). As a result, its implementation necessitated extreme violence to the theory of general relativity – something few physicists were prepared to sacrifice in favour of ‘numerology’ – requiring a time-dependent gravitational coupling of the form $G(t) \sim \frac{1}{t}$. Zel’dovich’s insight (1968) was to realise that a small but nonzero cosmological term ($\Lambda > 0$) allowed the present day radius of the Universe to be identified with the de Sitter radius, $r_U \sim \frac{1}{\sqrt{\Lambda}}$, which removed the need for time-dependence in the fundamental couplings. Thus, he obtained the formula $\Lambda \sim m^6 G^2 / \hbar^4$, where $m$ is a mass scale characterising the relative strengths of the gravitational and electromagnetic interactions, which he identified with the proton mass $m_p$. In this paper, we review a number of recent arguments which, instead, suggest the identification $m = m_e / \alpha$, where $m_e$ is the electron mass and $\alpha_e = e^2 / \hbar c \approx 1 / 137$ is the usual fine structure constant. We note that these are of a physical nature and, therefore, represent an attempt to lift previous arguments à la Dirac from the realm of numerology into the realm of empirical science. If valid, such arguments suggest an intimate connection, not only between the macroscopic and microscopic worlds, but, perhaps even more surprisingly, between the very essence of “dark” and “light” physics.

1. Introduction – Dirac, Zel’dovich and the Large Number Hypothesis

In 1937, Dirac noted the approximate order of magnitude equivalence between several large dimensionless numbers obtained from atomic physics and cosmology [1]. These included the ratio of the present day radius of the Universe, $r_U \approx 4.40 \times 10^{28}$ cm, to the classical electron radius, $r_e = 2.818 \times 10^{-15}$ cm, and the ratio of the electric and gravitational forces between an electron and a proton,

$$\frac{r_U}{r_e} \approx 10^{40} \approx \frac{e^2}{G m_e m_p} \approx 10^{39},$$

where $m_e = 9.109 \times 10^{-28}$ g and $m_p = 1.672 \times 10^{-24}$ g. (For convenience, from here on, we use the most recent measured, or inferred, values of physical length and mass scales.) Interpreting this as a signature of an as yet unknown connection between cosmological and atomic physics,
he formulated the Large Number Hypothesis (LNH), which states that the numerical equality between two very large quantities with similar physical meanings cannot be a simple coincidence [1, 2, 3, 4]. However, since the radius of the Universe is not constant but scales approximately as $r_U(t) \sim t$, Dirac surmised that the gravitational coupling varies according to $G(t) \sim 1/t$ — under the assumption $\Lambda = 0$ — though this is, of course, not the only way to satisfy Eq. (1) for all epochs [5]. Nonetheless, such variation is strongly at odds with the current ‘concordance’ model of cosmology, which is based on Einstein’s theory with a positive cosmological term and an unknown “dark” matter component [6].

In 1968, Zeldovich noted the same approximate equivalence between the ratio of the radius of the Universe and the Compton wavelength of the proton, $\lambda_p = h/m_p c = 1.321 \times 10^{-13}$ cm, and between $\lambda_p$ and the proton’s Schwarzschild radius, $r_S(m_p) = 2Gm_p/c^2 = 2.484 \times 10^{-52}$ cm [7]. In addition, he noted that, if $\Lambda > 0$ and $r_U \simeq l_{dS} \simeq 1/\sqrt{\Lambda}$ — contrary to Dirac’s original assumptions — then

$$\frac{r_U}{\lambda_p} \simeq \frac{m_p c}{h\sqrt{\Lambda}} \simeq 10^{40} \simeq \frac{\lambda_p}{r_S(m_p)} \simeq \frac{hc}{Gm_p^2} \simeq 10^{39},$$

and hence

$$\Lambda \simeq \frac{m_p^6 G^2}{h^4} \simeq 10^{-56} \text{ cm}^{-2}. \quad (3)$$

Setting $\Lambda \simeq 1/l_{dS}^2$, where $l_{dS}$ is the de Sitter horizon, this is equivalent to

$$\lambda_p = \frac{h}{m_p c} \simeq \left(\frac{r_{Pl}}{l_{dS}}\right)^{1/3}, \quad (4)$$

with $l_{dS} \simeq 10^{28}$ cm, where $r_{Pl} = \sqrt{hG/c^3}$ is the Planck length. Zel’dovich’s formula (3) is remarkable in that, not only does it establish a possible connection between dark energy ($\Lambda$), canonical gravity ($G$) and elementary particle physics ($m_p$), it is also compatible with current experimental bounds on the dark energy density [8, 9].

The current best fit to the available cosmological data favours a cosmological concordance, or $\Lambda$CDM model, in which dark energy takes the form of a positive cosmological constant, which accounts for approximately 69% of the total energy density of the Universe, whereas cold dark matter (CDM) accounts for around 26% and ordinary visible matter for around 5% [8, 9]. The present day density is close to the critical density,

$$\rho_{\text{crit}} = \frac{3H_0^2}{8\pi G} = 8.639 \times 10^{-30} \text{ g cm}^{-3}, \quad (5)$$

where $H_0 = 67.74$ km s$^{-1}$ Mpc$^{-1} = 2.198 \times 10^{-18}$ s$^{-1}$ is the Hubble constant, so that the dark energy density is

$$\rho_\Lambda = \frac{\Lambda c^2}{8\pi G} = 5.971 \times 10^{-30} \text{ g cm}^{-3}. \quad (6)$$

This yields an inferred value of the cosmological constant,

$$\Lambda = 1.114 \times 10^{-56} \text{ cm}^{-2}, \quad (7)$$

which corresponds to a de Sitter radius, $l_{dS} = 1.641 \times 10^{28}$ cm.
We now note the approximate equivalence between the classical electron radius and the proton Compton wavelength,

\[ \lambda_p = \frac{\hbar}{m_e c^2} = 1.321 \times 10^{-13} \text{cm} \]

\[ \simeq r_e = \frac{e^2}{m_e c^2} = 2.818 \times 10^{-13} \text{cm}, \quad (8) \]

or \( \lambda_p \simeq r_e = \alpha_e k^{-1}_G \), where \( k^{-1}_G = 2\pi/\lambda_e = \hbar/m_e c \) is the reduced Compton wavelength of the electron and \( \alpha_e = e^2/\hbar c = 7.927 \times 10^{-3} \) is the fine structure constant. Performing the replacement \( \lambda_p \rightarrow r_e \) in Eq. (3) then yields

\[ \Lambda \simeq \frac{m_6^6 G^2}{\alpha_6^6 \hbar^4} \simeq 1.366 \times 10^{-56} \text{cm}^{-2}, \quad (9) \]

or, equivalently,

\[ r_e = \frac{e^2}{m_e c^2} \simeq \left( \frac{l_{Pl}^2 dS}{r} \right)^{1/3}, \quad (10) \]

where \( l_{Pl} = \sqrt{\hbar G/c^3} \) denotes the reduced Planck length.

Remarkably, since the publication of Zel’dovich’s seminal paper [7], the relations (9)-(10) have been obtained in the literature using at least four independent, yet not inconsistent, methods. In the present work, we review these four main, existing approaches, and briefly discuss directions for future research. The outline of this paper is as follows. In Sec. 2, the four approaches are reviewed in chronological order (i.e., the order in which they were proposed in the literature), in subsections 2.1-2.4, respectively. In Sec. 3, the implications of the relations (9)-(10) for contemporary issues in theoretical physics, including holography [10, 11] and the present day accelerated expansion of the Universe [8, 9], are considered, along with potentially novel implications for the early Universe. A brief summary of our conclusions and main results is given in Sec. 4.

2. Beyond numerology – physical arguments for the LNH

2.1. Renormalization group flow of the vacuum energy

The relations (9)-(10) were first obtained in 1993 by Nottale [12] who argued that, like other fundamental ‘constants’, the cosmological constant is an explicitly scale-dependent quantity, obeying a renormalization group equation. As such, its present day value may be split into a ‘bare’ gravitational part and a scale-dependent, quantum mechanical vacuum energy part, i.e. \( \Lambda(r) = \Lambda_G + \Lambda_Q(r) \).

Following Zel’dovich [7], who also noted that the bare zero-point energy is unobservable, and who suggested that the observable contribution to the vacuum energy is given by the gravitational energy of virtual particle-antiparticle pairs, continually created and annihilated in the vacuum state,

\[ E_{grav} \simeq Gm^2(r) \frac{1}{r}, \quad (11) \]

where \( m(r) \simeq \hbar/(cr) \) is the effective mass of the particles at scale \( r \), Nottale obtained the scale-dependent formula for the vacuum energy density as

\[ \rho_{vac}(r) \simeq \rho_{Pl} \left( \frac{l_{Pl}}{r} \right)^6. \quad (12) \]
Here, \( \rho_{Pl} = (3/4\pi)m_{Pl}/l_{Pl}^3 \) is the Planck density, and \( m_{Pl} = \sqrt{\hbar c/G} \) denotes the reduced Planck mass.

Further assuming a renormalization group equation of the form

\[
\frac{d\rho_{\text{vac}}}{d\ln(r)} = \gamma(\rho_{\text{vac}}),
\]

where \( \gamma(\rho_{\text{vac}}) \) is an unknown function, which may be expanded as \( \gamma(\rho_{\text{vac}}) \approx \gamma_0 + \gamma_1 \rho_{\text{vac}} \) to first order, for \( \rho_{\text{vac}} \leq \rho_{Pl} \), yields

\[
\rho_{\text{vac}}(r) \approx \rho_0 \left[ 1 + \left( \frac{r_0}{r} \right)^{-\gamma_1} \right],
\]

where \( \rho_0 = -\gamma_1/\gamma_0 \) and \( r_0 \) is an integration constant. Comparison of Eqs. (13)-(14) then gives \( \gamma_1 = -6 \) and \( \rho_0 = \rho_{Pl} / (6/\rho_{Pl}) \).

Thus, Nottale obtained the low-energy asymptotic value of the cosmological constant, which was found to be scale-independent. Next, he argued that the transition between scale-dependence and scale-independence should be identified with the Thomson scattering length (the classical electron radius), given via \( \sigma_T \approx \pi r_e^2 \), where \( \sigma_T \) is the scattering cross-section. This is equal to the \( e^+e^- \) annihilation cross-section

\[
\sigma(e^+e^-) = \pi r_e^2 \left( \frac{m_e c^2}{E} \right),
\]

evaluated at \( E \approx m_e c^2 \).

In other words, \( r_e \) represents the effective electron radius, which is an energy-dependent quantity \( r(E) \), evaluated at its own mass scale. The cross-section for \( e^+e^- \) pair-production at this energy scale represents the main contribution to the vacuum energy at late times. Hence, by identifying \( \rho_{\text{vac}} \equiv \rho_\Lambda \) and \( r_0 \equiv r_e \) in Eq. (11), he obtained the relation

\[
\Lambda \approx \frac{\rho_{Pl}}{r_e^6},
\]

which is equivalent to Eqs. (9)-(10).

2.2. Dark energy particles and the ‘Small Number Hypothesis’

In 2008, Eqs. (9)-(10) were also obtained by Boehmer and Harko [13] in 2008, who noticed that the length scale given by the right-hand side of Eq. (4) represents the largest stable radius of a particle of mass

\[
m_{dS} \equiv \frac{\hbar}{l_{dS} c} \approx 10^{-66}\text{g},
\]

and density \( \rho_{\text{min}} \approx \rho_\Lambda \). The mass scale \( m_{dS} \) was previously proposed by Wesson as the minimum possible mass/energy scale in the present day Universe [14], and the existence of a minimum possible density, \( \rho_{\text{min}} = \rho_\Lambda / 2 \), for stable, charge-neutral, self-gravitating compact objects in the presence of a positive cosmological constant (\( \Lambda > 0 \)), was shown in [15]. This result was obtained directly from the generalised Buchdahl inequalities

\[
\frac{2Gm}{c^2} \geq \frac{\Lambda}{6} R^3, \quad \rho = \frac{3m}{4\pi R^3} \geq \rho_{\text{min}} \equiv \frac{\Lambda c^2}{16\pi G},
\]
where \( R \) represents the classical radius of the compact object. Substituting \( m_{\text{dS}} \) into Eq. (18), we obtain

\[
R \leq (t_p^2 m_{\text{dS}})^{1/3} \approx 10^{-13} \text{cm}.
\]

Interestingly, it may be shown that this length scale also represents the minimum radius into which the present mass of the Universe, \( m_U \approx m'_{\text{dS}} \equiv m_{\text{Pl}}^2 / m_{\text{dS}} \), can be compressed without exceeding the Planck density [16].

Noting the numerical coincidence between this length scale and the classical electron radius, Boehmer and Harko proposed a formal equivalence between the two on the basis of a ‘Small Number Hypothesis’ (SNH), which directly yields Eq. (16). However, since the reciprocal of a large number is a small number, this may be considered logically equivalent to Zel’dovich’s ‘Number Hypothesis’ (SNH), which directly yields Eq. (16). However, since the reciprocal of a large number is a small number, this may be considered logically equivalent to Zel’dovich’s reformulation of the LNH for \( \Lambda > 0 \): what differs is the empirical content, according to the substitution \( \lambda_{\text{Pl}} \to r_e \). Thus, the work presented in [13, 16] shows that, not only is the length scale given in Eq. (19) comparable to the classical electron radius, it also represents (i) the maximum radius of a minimum-mass, minimum-density particle and (ii) the minimum radius of a maximum-mass, maximum-density particle in the observable Universe.

However, for charge-neutral, quantum mechanical particles, we may identify the radius \( R \) with the Compton wavelength \( \lambda_C = \hbar / (mc) \) (from here on, we use conventional notation and terminology, referring to the Compton wavelength and reduced Compton wavelength interchangeably), yielding [16]

\[
m \geq m_{\Lambda} \equiv \frac{1}{\sqrt{2}} \sqrt{m_{\text{Pl}} m_{\text{dS}}} \approx 10^{-35} \text{g}.
\]

Hence, we see that \( E_{\text{dS}} = m_{\text{dSc}}^2 \) may be interpreted as the minimum possible quantum of energy – corresponding to a de Broglie wavelength of the order of the de Sitter horizon – but cannot be the minimum rest mass of a stable massive particle.

The mass scale \( m_{\Lambda} \) (20) has several interesting properties. According to the model presented in [16], it represents the minimum possible mass of a stable, charge-neutral, quantum mechanical and self-gravitating body in the presence of dark energy (\( \Lambda > 0 \)). With this in mind, it is notable that it is consistent with current experimental bounds on the mass of the electron neutrino obtained from Planck satellite data [9]. It may also be interpreted as the mass of an effective dark energy particle and its associated Compton wavelength, \( \ell_\Lambda = \sqrt{2} / t_p m_{\text{dS}} \), is of the order of 0.1 mm. Thus, according to this model, the dark energy density is approximately constant over large distances, but may become granular on sub-millimetre scales. In this context, it is notable that that tentative hints of periodic variation in the gravitational field strength on precisely this length scale have recently been observed [17]. It is also the unique mass scale for which the Compton wavelength is equal to the gravitational turn-around radius in the presence of \( \Lambda > 0 \) [18].

As we shall see in Sec. 2.4, \( m \approx m_{\Lambda} \) also arises naturally in the context of the the dark-energy modified uncertainty principle (DE-UP), proposed in [19, 20], together with the mass scale \( m \approx m_e \approx \alpha_e (m_{\text{Pl}} m_{\text{dS}})^{1/3} \). The former corresponds to an absolute minimum, while the latter is associated with the length scale (19) via the relativistic formula for the minimum stable radius of a charged particle, \( R \geq Q^2 / (mc^2) \), evaluated for \( Q = \pm e \).

2.3. Information theory

Also in 2008, Eqs. (9)-(10) were obtained by Beck using an information theoretic approach to the cosmological constant problem [21]. He used a system of four axioms, constructed by analogy with the Kinchin axioms of information theory [22], which describe the most ‘desirable’
properties of an information measure, to fix the form of the cosmological constant in terms of the other fundamental constants of nature.

It may be shown that the Kinchin axioms uniquely fix the form of the Shannon entropy, which forms the mathematical, though not the microphysical, basis of statistical mechanics and thermodynamics. Thus, by constructing an analogous approach to the cosmological constant problem, Beck attempted to fix the form of $\Lambda$ in terms of the other constants of nature on an axiomatic basis, without reference to an underlying microphysical theory.

By formally replacing the dependence of the information measure $I$ on the probabilities of events $p_i$ by the dependence of $\Lambda$ on the remaining physical constants, i.e., the fundamental coupling constants $\alpha_i$, masses $m_i$ and mixing angles $s_i$, he argued that the requirements of Fundamentality (L1), Boundedness (L2), Simplicity (L3) and Scale-invariance (L4) uniquely fix the form of the cosmological constant according to Eqs. (9)-(10). These axioms are constructed by analogy with the four Kinchin axioms (K1-K4) of the same names.

Specifically, K1, ‘fundamentality’, simply states that the information measure $I$ should depend on the fundamental quantities, the probabilities of events $p_i$, 

$$I = I(p_i), \quad (K1)$$

and not on any other factors. The analogous axiom L1 for the cosmological constant is 

$$\Lambda = \Lambda(\{\alpha_i\}, \{m_i\}, \{s_i\}), \quad (L1)$$

K2, ‘boundedness’ states that there exists a lower bound for the value of $I$, corresponding to the uniform distribution, $p_i = 1/N$, where $N$ is the total number of distinct events, such that 

$$I(1/N, 1/N...1/N) \leq I(p_1, p_2...p_N), \quad (K2)$$

The analogous axiom L2 states 

$$0 < \Lambda, \quad (L2)$$

where $\Lambda = 0$ is explicitly excluded. K3, ‘simplicity’, states that the information measure should not change if the set of events is enlarged by another set with probability zero, i.e. 

$$I(p_1, p_2...p_N) = I(p_1, p_2...p_N; 0), \quad (K3).$$

The analogous axiom L3 is 

$$\Lambda(\{\alpha_i\}, \{m_i\}, \{s_i\}) = \Lambda(\{\alpha_i\}, \{m_i\}, \{s_i\}; \{c_i\}), \quad (L3),$$

where the $c_i$ are not fundamental constants of nature. The final axiom, K4 ‘invariance’ is the most restrictive. It may be expressed as 

$$I(\{p^{\text{I}}_{ij}\}) = I(\{p^\text{I}_i\}) + \sum_i p^\text{I}_i I(\{p^\text{II}(j|i)\}), \quad (K4)$$

where the superscripts $I$ and $\text{II}$ denote probabilities of events in different (not necessarily independent) subsystems and $p^\text{II}(j|i)$ is the conditional probability of event $j$ in subsystem $\text{II}$, given an event $i$ in subsystem $I$. $I(\{p^\text{II}(j|i)\})$ is the conditional information of subsystem $\text{II}$, in the joint system $I, \text{II}$, described by the probabilities $p^\text{I}\text{II}_{ij} = p^\text{I}_i p^\text{II}(j|i)$. The meaning of K4 is that the information measure should be independent of the way in which the information is collected. We may either (a) collect information in subsystem $I$, then in subsystem $\text{II}$, or (b)
collect information in subsystem II, assuming a given event in subsystem I, before summing over all possible events in subsystem I, weighted by their respective probabilities $p_i$. Hence, there is a scale transformation in the space of possible information measures and probabilities, such that

$$I(\{\tilde{p}_i\}) = \tilde{I}(\{p_i\}), \quad (K4),$$

where a tilde denotes transformed quantities. The analogous axiom L4 may therefore be expressed as

$$\Lambda(\{\tilde{\alpha}_i\}, \{\tilde{m}_i\}, \{\tilde{s}_i\}) = \tilde{\Lambda}(\{\alpha_i\}, \{m_i\}, \{s_i\}), \quad (K4),$$

where a tilde also denotes an appropriate scale transformation: in this case, a literal rescaling of the fundamental constants of nature by an arbitrary numerical factor.

After formulating the axioms L1-L4, the argument presented in [21] relies on three specific assumptions, (i) gravitational scale invariance, (ii) a dimensional argument and (iii) electromagnetic scale invariance. The first and last of these are implicit in axiom L4, but Beck singled out the electromagnetic and gravitational interactions as being of greatest relevance to the large-scale dynamics of the early (before recombination) and late-time (after recombination) Universe, respectively. We now consider each of these assumptions in detail.

(i) Gravitational scale-invariance: In the Newtonian approximation, the the gravitational energy density of a distribution of point-like masses occupying a volume $V$ is

$$\rho_G = -\frac{G}{V} \sum_{i,j} \frac{m_i m_j}{r_{ij}}, \quad (30)$$

where $r_{ij}$ denotes the distance between the $i^{th}$ and $j^{th}$ masses. Thus, if the gravitational constant $G$ is rescaled such that

$$G \rightarrow \Gamma G, \quad (31)$$

where $\Gamma$ is an arbitrary numerical constant, but the masses $m_i$, $m_j$ and distances $r_{ij}$ are kept the same, the energy density scales as

$$\rho_G \rightarrow \Gamma \rho_G. \quad (32)$$

Scale invariance of the ratio $\rho_{\text{vac}}/\rho_G$ then requires

$$\rho_{\text{vac}} \rightarrow \Gamma \rho_{\text{vac}}. \quad (33)$$

Hence, $\rho_{\text{vac}} \propto G$ and we may set

$$\rho_{\text{vac}} \sim G X, \quad (34)$$

where $X$ is an, as yet unknown, quantity.

(ii) The dimensional argument: On purely dimensional grounds, the unknown factor $X$ must take the form $X \sim (c^4/\hbar^4)m^6$, where $m$ is an arbitrary mass scale, composed (in some way) from fundamental mass scales, dimensionless coupling constants and mixing angles, according to axiom L1. (Note that this dimensional argument is unaffected even if the fundamental constants that are not themselves fundamental couplings, masses or mixing angles, i.e. $\hbar$ and $c$, are also rescaled by $\Gamma$..) Thus, without loss of generality, we may set

$$\rho_{\text{vac}} = A \frac{c^4}{\hbar^4} m^6, \quad (35)$$

where
where $A$ is a numerical constant, which may depend on the remaining constants of nature \{\{\alpha_i\}, \{m_i\}, \{s_i\}\} in any way. Likewise, $A$ may be expressed arbitrarily in terms of $\alpha_e$, such that

$$\rho_{\text{vac}} = \frac{1}{8\pi} \alpha_e^\eta \frac{c^4}{\hbar^4} m_e^6,$$

where $\eta$ is an arbitrary function of \{\{\alpha_i\}, \{m_i\}, \{s_i\}\}.

(iii) Electromagnetic scale-invariance: Beck then argued that the most important process determining the large-scale evolution of the Universe before recombination was the Thomson scattering of electrons, associated with the scattering length $r_e \sim \sqrt{\sigma_T}$. Clearly, this line of reasoning is similar to Nottale’s and Beck could also have argued à la Nottale that $r_e$, interpreted as the effective cross-sectional radius for $e^+e^-$ pair-production, is just as relevant to the large-scale dynamics of the late-time Universe. Finally, therefore, Beck claimed that, since simultaneous scale transformations of the form

$$\alpha_e \rightarrow \Gamma \alpha_e, \quad m_e \rightarrow \Gamma m_e,$$

leave the Thomson scattering cross-section $\sigma_T$ invariant, scale-invariance of the ratio $\Lambda^{-1}/\sigma_T$ requires that $\Lambda$ be dependent on the ratio $m_e/\alpha_e$ only. This fixes the remaining free constant, $\eta = -6$, and the factor of $(8\pi)^{-1}$ in Eq. (36) is chosen so as to match the numerical factor in Eq. (6), giving the ‘simplest’ result:

$$\rho_\Lambda = \frac{\Lambda c^2}{8\pi G} \equiv \rho_{\text{vac}} = \frac{1}{8\pi} \frac{c^4}{\hbar^4} \left( \frac{m_e}{\alpha_e} \right)^6.$$

It is straightforward to check that this is relation is equivalent to Eqs. (9)-(10) and (16).

However, it must be noted that an obvious problem with Beck’s approach is that, unlike the probabilities of events, which are all dimensionless, the fundamental quantities on which he bases the expression for the cosmological constant $\Lambda$ are inequivalent in this respect. Though non-gravitational coupling constants may be expressed in dimensionless form using $\hbar$ and $c$, Newton’s constant $G$ cannot. Similarly, the fundamental masses cannot be expressed in dimensionless form without the use of the Planck mass, $m_{\text{Pl}} \propto 1/\sqrt{G}$. A related problem concerns the fact that, if the ultimate origin of the point-like masses considered in Eq. (30) are fundamental masses, these should also be rescaled at the same time as $G$, which changes the scaling of $\rho_{\text{vac}}$ with $\Gamma$, dramatically.

Overall, one can say that it does not make physical sense to rescale dissimilar quantities (i.e., those with inequivalent dimensions) with the same scale factor: only similar quantities should be rescaled in this way. A self-consistent set of rescalings, which leave dimensionless quantities constructed from Eq. (16) explicitly invariant (in line with axiom L4), would then be

$$\Lambda^{-1} \rightarrow \Gamma \Lambda^{-1}, \quad \sigma_T \rightarrow \Gamma \sigma_T, \quad l_\text{Pl}^2 \rightarrow \Gamma l_\text{Pl}^2.$$

Alternatively, instead of requiring the absolute scale independence of $\Lambda$, or even the relative scale independence suggested by Eq. (39), we may instead require holographic relationships between the bulk and the boundary of the Universe to be preserved by our expression for $\Lambda$, written in terms of the remaining constants of nature. In the following section, we summarize the results of a new approach to deriving the expressions (9)-(10) and (16), which automatically implements the holographic principle [10, 11].
2.4. Minimum length uncertainty relations in a dark energy Universe

A further attempt to explain the connection between cosmological and atomic scales in the LNH, in terms of the stability of fundamental particles in the presence of dark energy, was presented in a series of recent papers [19, 20]. (See also [23, 24] for extensions to modified gravity theories.) This approach considered the status of minimum length uncertainty relations (MLURs), motivated by quantum gravity phenomenology (see [25, 26] for contemporary reviews), in the presence of a vacuum energy density $\rho_{\Lambda}$ given by Eq. (6).

In [19], a new form of MLUR, dubbed the ‘dark energy uncertainty principle’ or DE-UP for short, which explicitly includes the de Sitter scale $l_{dS} \sim 1/\sqrt{\Lambda}$, as well as the Planck scale, was proposed. The general form of the DE-UP proposed in [19] is

$$\Delta x_{\text{total}}(\Delta v, r, m) = \Delta x_{\text{canon}}(\Delta v, r, m) + \Delta x_{\text{grav}}(r, m)$$

$$\geq \Delta x(\Delta v) + \Delta x_{\text{recoil}}(\Delta v, r, m)$$

$$+ \Delta x_{\text{grav}}(r, m)$$

$$\geq (\Delta x_{\text{canon}})_{\text{min}}(r, m)$$

$$+ \Delta x_{\text{grav}}(r, m). $$

Here, $(\Delta x_{\text{canon}})_{\text{min}}$ denotes the minimum possible canonical quantum uncertainty of a wave packet that has been freely evolving for a time $t = r/c$. By explicitly solving the free-particle Schrödinger equation in the Heisenberg picture, it is straightforward to show that, for a particle of mass $m$, the minimum canonical uncertainty is

$$(\Delta x_{\text{canon}})_{\text{min}} \simeq \sqrt{\lambda_C r}, \quad (41)$$

where $r = ct$ and $\lambda_C$ is the Compton wavelength [27, 28]. This expression can also be obtained by considering a gedanken experiment, originally due to Salecker and Wigner, in which a massive particle is used to ‘measure’ a distance $r \approx ct$ by means of the emission and reabsorption of a photon [29]. By minimising the sum of the first two terms on the top line of Eq. (40), where $\Delta x$ represents the canonical Heisenberg uncertainty and $\Delta x_{\text{recoil}} \simeq \Delta vt \equiv \Delta pr/(mc)$ is the additional uncertainty due to recoil, with respect to either $\Delta v$ or $m$, we obtain Eq. (41) directly.

The term $\Delta x_{\text{grav}}$ represents an additional contribution to the total uncertainty, due to the superposition of gravitational field states which are (in turn) induced by the superposition of position states associated with $m$. From a relativistic perspective, this may be considered equivalent to the uncertainty associated with the superposition of space-time geometries in the quantum gravity regime. In [19], it was conjectured that this superposition is influenced by two factors: (i) the mass of the particle, and (ii) the presence of the dark energy density, or, equivalently, of a finite horizon for the wave function centre of mass, $l_{dS} \sim 1/\sqrt{\Lambda}$. Taking both these factors into account, the conjectured form of $\Delta x_{\text{grav}}$ proposed in [19] was

$$\Delta x_{\text{grav}} \simeq \frac{G m}{\sqrt{\Lambda} c^2 r} \simeq \frac{\ell_{Pl}^2 l_{dS}}{\lambda_C r}. \quad (42)$$

This was combined with the limit on the mass/radius ratio for stable, charged, compact objects originally obtained by Bekenstein [30], and later generalised by Boehmer and Harko for $\Lambda > 0$ [31], i.e.

$$R \geq \frac{Q^2}{mc^2} + \text{h.o.t.}(G, \Lambda, \ldots). \quad (43)$$

To leading order, Eq. (43) simply recovers the well-known formula for the classical radius of a charged particle obtained from special relativity, but this remains rigorously valid in the weak-field limit of general relativity, even in the presence of dark energy [30, 31]. By identifying the
total minimum quantum mechanical uncertainty \((\Delta x_{\text{total}})_{\text{min}}\), including both canonical quantum and gravitational/dark energy-induced terms, with the minimum radius \(R\) obtained from the Bekenstein bound, the following inequality is obtained

\[ \frac{Q^2}{mc^2} \leq (l_{\text{Pl}}^2 l_{\text{dS}})^{1/3}. \]  

(44)

Thus, according to the model presented in [19], Eq. (44) gives the maximum possible charge-squared to mass ratio for a stable, charged, self-gravitating and quantum mechanical object. Assuming saturation of this bound is equivalent to assuming the existence of a particle in nature that simultaneously saturates both the classical (general-relativistic) and quantum mechanical (MLUR) stability constraints. Remarkably, comparison with Eq. (10) reveals that saturation is obtained for the electron charge-squared to mass ratio. Equivalently, we see that evaluating (44) for \(Q = \pm e\), yields

\[ m \geq \alpha_e (m_{\text{Pl}} m_{\text{dS}})^{1/3} = 7.332 \times 10^{-28} \text{ g} \]

\[ \simeq m_e = 9.109 \times 10^{-28} \text{ g}. \]  

(45)

Hence, if the electron were more highly charged (with the same mass \(m_e\)) or any less massive (with the same charge \(e\)), a combination of electrostatic and dark energy repulsion would destabilize its Compton wavelength [19].

In [20], a full physical derivation of the dark energy-modified uncertainty principle (DE-UP), proposed in [19] is given, in which it is shown that the de Sitter length-dependent term arises as a direct consequence of the existence of a finite horizon \(r_H(\tau_0) \approx l_{\text{dS}}\), where \(\tau_0\) is the present age of the Universe. In addition, it is shown that the DE-UP defined by Eqs. (40) and (41)-(42) naturally incorporates the mass bound for neutral particles, Eq. (20), in addition to that for charged particles, Eq. (45). In particular, it is straightforward to show that requiring every potentially observable length-scale, i.e. \((\Delta x_{\text{canon}})_{\text{min}}, \Delta x_{\text{grav}}\) and the ‘probe’ distance \(r\), to be super-Planckian, automatically implies the existence of a minimum mass in nature, \(m_\Lambda \simeq \sqrt{m_{\text{Pl}} m_{\text{dS}}}\). Alternatively, beginning with the bound (20), obtained by combining both general-relativistic and quantum mechanical effects, we obtain \(l_{\text{Pl}}\) as a limiting resolution for physical measurements of length within the DE-UP framework.

Finally, we note that, using Salam’s theory of strong gravity [32, 33, 34, 35] as an effective theory for modeling quark confinement, analogous arguments where applied to charged, strongly interacting particles, in which the energy density associated with the ‘strong de Sitter radius’ was identified with the ‘bag constant’ of the MIT bag model for confined nuclear matter, \(B \simeq 10^{14} \text{ gc当地^{-1}}\) [24]. These arguments successfully predicted the correct order of magnitude value for the charge-squared to mass ratio of the up quark, the lightest known strongly interacting and quantum mechanical particle, with charge \(Q = \pm 2e/3\) [24].

3. Implications of a connection between “dark” and “light” physics

3.1. Holography

It is straightforward to see that, for any particle that minimizes the total uncertainty given by the DE-UP, Eqs. (40) and (41)-(42), a holographic relation holds between the bulk and the boundary of the Universe, namely

\[ \left( \frac{(\Delta x_{\text{total}})_{\text{min}}}{l_{\text{dS}}} \right)^3 = \frac{l_{\text{Pl}}^2}{l_{\text{dS}}^2} = N = 1.030 \times 10^{122}. \]  

(46)

Hence, the number of Planck sized ‘bits’ on the de Sitter boundary is equal to the number of minimum-volume ‘cells’, \(V_{\text{cell}} \sim (\Delta x_{\text{total}})_{\text{min}}^3\), in the bulk [19]. As mentioned previously, in Sec.
2.2, R \simeq (\Delta x_{\text{total}})_{\text{min}} may also be interpreted as the classical radius of a ‘particle’ with both minimum energy, \(E_{\text{dS}} = m_{\text{dS}}c^2\), and minimum density, \(\rho_{\text{min}} = \rho_{\Lambda}/2\), i.e.

\[
\rho \simeq \frac{m_{\text{dS}}}{R^3} \geq \rho_{\text{min}} \simeq \frac{\Lambda c^2}{G} \iff R \leq (\Delta x_{\text{total}})_{\text{min}} \simeq (l_{\text{Pl}}^2 l_{\text{dS}})^{1/3}.
\]

Although a massive particle with rest energy \(E_{\text{dS}}\) would be unstable due to the repulsive effects of dark energy, this may correspond to the energy of a photon with maximum wavelength, \(\lambda \simeq l_{\text{dS}}\). Thus, \((\Delta x_{\text{total}})_{\text{min}}\) may also be interpreted as the classical radius of a localized, minimum-energy photon. A space-filling ‘sea’ of such photons would have the same energy density as the dark energy field [16].

In addition, we may consider a maximum-mass, maximum-density state, for which \(\rho \simeq \rho_{\text{Pl}}\) and the total energy is \(E'_{\text{dS}} \simeq m'_{\text{dS}}c^2\) \((m'_{\text{dS}} \equiv m_{\text{Pl}}/m_{\text{dS}})\). The classical radius thus obtained corresponds to the smallest possible volume within which the total mass of the present day horizon may be confined, without exceeding the Planck density. We then have [19]

\[
\rho \simeq \frac{m'_{\text{dS}}}{R^3} \leq \rho_{\text{Pl}} \simeq \frac{c^5}{\hbar G^2} \iff R \geq (\Delta x_{\text{total}})_{\text{min}} \simeq (l_{\text{Pl}}^2 l_{\text{dS}})^{1/3}.
\]

The length scale \((\Delta x_{\text{total}})_{\text{min}} \simeq (l_{\text{Pl}}^2 l_{\text{dS}})^{1/3}\) therefore corresponds to at least three interesting physical scenarios. It may be interpreted as (i) the classical radius of a ‘particle’ with both minimum energy and minimum energy density, (ii) the classical radius of a ‘particle’ with both maximum energy and maximum energy density, and (iii) the classical radius/minimum total uncertainty of the electron, which saturates the charged particle stability bound (44). In the context of the DE-UP model [19, 20], all three interpretations satisfy the general holographic relation, Eq. (46).

3.2. Universal expansion

As pointed out in [16], the minimum mass \(m_{\Lambda} \simeq \sqrt{m_{\text{Pl}} m_{\text{dS}}}\) (20) may also be interpreted as the mass of an effective dark energy particle. Since, in this model, even random quantum fluctuations reduce the inter-particle distance between nearest neighbours to less than \(\lambda_{\text{C}}(m_{\Lambda}) = l_{\Lambda} \simeq l_{\text{Pl}} l_{\text{dS}}\), it follows that the pair-production of “dark” minimum-mass particles is capable of driving the present day accelerated expansion of the Universe. In short, if space is ‘full’ of dark energy particles, with mean inter-particle distance \(l_{\Lambda} \sim 0.1\) mm, the pair-production necessitated by quantum mechanics requires a concomitant expansion of space [16, 24].

Let us assume that the probability of a single (holographic) spatial cell ‘pair-producing’ within a time interval \(\Delta \tau = t_{\text{Pl}} = l_{\text{Pl}}/c\), due to the pair-production of dark energy particles, is given by

\[
P(\Delta V = +V_{\text{cell}}|V_0 = V_{\text{cell}}, \Delta \tau = t_{\text{Pl}}) \simeq N^{-1/2}
\]

\[
= V_{\text{Pl}} = \frac{l_{\text{Pl}}}{l_{\text{dS}}} \simeq (\frac{\hbar G \Lambda}{3 c^4})^{1/2} \simeq 9.851 \times 10^{-62},
\]

where \(V_0\) denotes the initial volume at the initial time. This leads naturally to a de Sitter-type expansion, modelled by the differential equation

\[
\frac{da^3}{d\tau} \simeq \frac{N^{-1/2} a^3}{t_{\text{Pl}}} = \frac{l_{\text{Pl}} a^3}{l_{\text{dS}} t_{\text{Pl}}},
\]

or, equivalently,

\[
\frac{da}{d\tau} \simeq c \sqrt{\frac{\Lambda}{3}} a, \quad a(\tau) \simeq a_0 e^{-c\sqrt{\Lambda/3} \tau}.
\]
Since the production of a single dark energy energy particle requires the production of \( n_{\text{cell}} = V_{\Lambda} / V_{\text{cell}} \approx l_{\Lambda}^3 / (l_{p}^2 l_{\text{ds}}) = N^{1/4} \) cells of space, this implies that the probability of a dark energy particle pair-producing within a single Planck time is given by

\[
P(\Delta V) = -V_{\Lambda} |V_0 = V_{\Lambda}, \Delta t = t_{p}) \approx N^{-3/4} \]

\[
= \left( \frac{l_{\text{ds}}}{l_{p}} \right)^{-3/2} \approx 10^{-91}. \tag{52}
\]

However, since there are \( n_{DE} \approx l_{\Lambda}^3 / l_{p}^3 = N^{3/4} \) dark energy particles within the de Sitter horizon, this means that one dark energy particle is produced, somewhere in the observable Universe, during every Planck time interval. This rate of pair-production is capable of giving rise to the accelerated Universal expansion observed at the present epoch.

In this model, the observed vacuum energy is really the energy associated with the dark energy field, for which \( \lambda_C(m_{\Lambda}) = l_{\Lambda} \) provides a natural a UV cut-off for the field modes, yielding

\[
\rho_{\text{vac}} \approx \frac{\hbar}{c} \int_{l_{\Lambda}}^{l_{\text{ds}}} \sqrt{k^2 + \left( \frac{2\pi}{l_{\Lambda}} \right)^2 - k^2} \frac{d^3k}{l_{\text{ds}}} \approx \frac{m_{\text{pl}} l_{p}}{l_{\Lambda}^3} \approx \frac{\Lambda_c^2}{G} \approx 10^{90} \text{ g cm}^{-3}. \tag{53}
\]

The field itself remains ‘trapped’ in a Hagedorn-type phase, in which any increase in kinetic energy, even that caused by random collisions between neighbouring dark energy particles due to quantum uncertainty, results in pair-production, rather than an increase in temperature.

The temperature associated with the field therefore remains constant, on large scales, and is comparable to the present day temperature of the CMB,

\[
T_{\Lambda} \equiv \frac{m_{\Lambda} c^2}{8\pi k_{\text{B}}} \approx 2.27 \text{ K } \approx T_{\text{CMB}} = 2.73 \text{ K}. \tag{54}
\]

Here the factor of \( 8\pi \) is included by analogy with the expression for the Hawking temperature,

\[
T_H \equiv \frac{c^2}{8\pi k_{\text{B}}} \frac{m_{\text{pl}}^2}{m}, \tag{55}
\]

so that \( T_{\Lambda} = T(m_{\Lambda}) = T_{\text{H}}(m_{\Lambda}') \), where \( m_{\Lambda}' \equiv m_{\text{pl}}^2 / m_{\Lambda} \) is the dual mass. Though this may seem like another ‘miraculous’ coincidence, à la Dirac, in the dark energy model implied by the DE-UP it is simply a re-statement of the standard ‘coincidence problem’ of cosmology, i.e. the Universe begins a phase of accelerated expansion when \( r_U \approx l_{\text{ds}} \), at which point \( \Omega_M \approx \Omega_{\Lambda} \) and, hence, \( T_{\text{CMB}} \approx T_{\Lambda} \). The question remains, why do we live at precisely this epoch? However, no new coincidences are required, in order to ‘explain’ Eq. (54).

Again, we may apply analogous arguments to strongly interacting particles by using Salam’s theory of strong gravity as an effective theory for confined (and deconfined) nuclear matter. These predict a genuine Hagedorn temperature for the quark-gluon plasma of order \( T_{\text{Hag}} \approx N^{1/10} T_{\Lambda} \approx 10^{12} \text{ K} \) [24].

### 3.3. Additional cosmological implications

The holographic relation Eq. (46) remains valid for all epochs, prior to the present day, under the substitution \( l_{\text{ds}} \rightarrow r_H(\tau) \), where \( r_H(\tau) \) is the physical horizon radius at cosmic time \( \tau \). However, under these circumstances, we note that the DE-UP model naturally implies both a time-dependent minimum mass for neutral particles and a time-dependent maximum charge-squared to mass ratio for charged particles, i.e.

\[
m_{\nu}(\tau) \geq \sqrt{m_{\text{pl}} m_{H}(\tau)}, \tag{56}
\]
where $m_H(\tau) \equiv \hbar/(r_H(\tau)c)$, and

$$\frac{e^2}{m_e}(\tau) \leq c^2(l_{Pl}^2r_H(\tau))^{1/3}, \quad (57)$$

respectively [20]. Equation (57) corresponds to a three-dimensional, time-dependent, holographic cell radius

$$(\Delta x_{\text{total}})_{\text{min}}(\tau) \simeq (l_{Pl}^2r_H(\tau))^{1/3}. \quad (58)$$

This is similar to the MLUR for an expanding Universe suggested by Ng [36], but with the cosmological horizon $r_H(\tau)$ in place of the Hubble horizon $H(\tau)/c$.

Though highly speculative, these relations imply an interesting form of ‘unification’ in the early Universe, with all masses tending to the Planck mass and all charges tending to the Planck charge, $q_{Pl} = \sqrt{\hbar c}$, as $r_H(\tau) \to l_{Pl}$. Alternatively, if the Planck-density threshold limits the radius of the ‘initial’ big bang horizon such that $r_H(\tau) \geq (l_{Pl}^2d_{S})^{1/3}$, as suggested by the results obtained in [19], the minimum holographic cell radius will be of order

$$R_{\text{min}} \simeq (l_{Pl}^2d_{S})^{1/9} \simeq 10^{-26} \text{ cm}. \quad (59)$$

Finally, we note that the length and time scales associated with the mass $m_\nu(\tau)$ (56), i.e.

$$r_\nu(\tau) = c t_\nu(\tau) \equiv \frac{\hbar}{m_\nu(\tau)c}, \quad (60)$$

also satisfy the ‘four-dimensional’ holographic relation

$$\left(\frac{r_H(\tau)}{r_\nu(\tau)}\right)^4 = \frac{r_H^2(\tau)}{l_{Pl}^2} = N(\tau) \leq 1.030 \times 10^{122}. \quad (61)$$

Hence, the DE-UP model strongly suggests time-variation of either, or both, $e$ and $m_e$, assuming that $\{G, c, \hbar, \Lambda\}$ are genuine universal constants. Similar arguments apply to the mass of the lightest neutral particle, previously identified with the mass of the electron neutrino, $m_\nu$.

For models involving temporal and/or spatial variation of fundamental constants, the situation is even more complicated, and it may be extremely difficult, in practice, to distinguish variation in $e$ and/or $m_e$, and $m_\nu$, from other effects. Particular classes of models in which the variation of physical constants should, automatically, imply a modification of the DE-UP formulae, Eqs. (40) and (41)-(42), include those with a running gravitational coupling [37, 38, 39, 40], variable speed of light [41, 42, 43, 44, 45, 46, 47, 48], or dynamical dark energy field [49, 50, 51, 52]. (See also [53, 54, 55, 56, 57] for current bounds on varying $\alpha_e$ theories, including their effects on cosmic string phenomenology [58, 59] and [60, 61, 62, 63, 64] for more general models involving variations of multiple physical constants.)

In fact, several models incorporating non-minimal couplings between dark energy and the electromagnetic sector have already been proposed in the literature, as solutions to problems in contemporary cosmology [65, 66, 67, 68, 69]. Although a thorough analysis of the cosmological implications of the DE-UP has not yet been attempted, the cosmological implications of $\Lambda \propto \alpha_e^{-6}$ cosmology were investigated in [70, 71], in the context of a time-varying fine structure constant. Nonetheless, its unusual predictions suggest that future observations and/or analysis of currently available data may be capable of falsifying the model [20].
4. Discussion

We have shown that the relations (9)-(10) and (16), which are equivalent to Zel’dovich’s reformulation of Dirac’s Large Number Hypothesis for a Universe with $\Lambda > 0$, under the identification $m \simeq m_e/\alpha_e$, are well motivated from a number of theoretical perspectives. Each of these goes beyond numerology and aims to base the seemingly incredible coincidences noted by Dirac, Zel’dovich and others on firm physical arguments. Specifically, we have outlined four sets of independent, yet not necessarily incompatible arguments, given in the literature, which give rise to the relation $\Lambda \simeq l_{Pl}^3/r_e^6$ (16), where $r_e = e^2/(m_e c^2)$ is the classical electron radius.

The first, proposed by Nottale [12] (1993), is based on the assumption that the vacuum energy density is dominated by the gravitational energy associated with $e^+e^-$ pair production. By perturbatively expanding the renormalisation group equation for $\rho_{vac}$, and identifying the transition from running to scale-independence with the energy scale $E \simeq m_e c^2$, he obtained Eq. (16) directly.

The second, proposed by Boehmer and Harko [13] (2008), identifies the maximum stable radius of a minimum-mass, minimum-density ‘particle’, in the presence of a positive cosmological constant $\Lambda > 0$, with the classical electron radius via a ‘Small Number Hypothesis’. Since the reciprocal of a large number is a small number, this is logically equivalent to Zel’dovich’s reformulation of the LNH with $m \simeq m_e/\alpha_e$. However, the important physical content of this work is the realisation that the cosmological constant automatically implies the existence of a minimum density for stable compact objects in nature. This result arises rigorously from the generalized Buchdahl inequalities [15] and, following a similar analysis, the minimum charge-squared to mass ratio for electrically charged particles can also be obtained. Identifying this with the alternative expression for $r_e$, obtained from the SNH, also yields (16).

The third method, proposed by Beck [21] (2008) follows an axiomatic approach, based on analogy with the Kinchin axioms of information theory [22]. These uniquely fix the form of the Shannon entropy, which forms the mathematical basis of statistical mechanics and thermodynamics, without making any assumptions about the underlying microphysical basis of these theories. Likewise, Beck’s approach aims to uniquely fix the form of $\Lambda$ without reference to an underlying microphysical model. By formally replacing the dependence of the Shannon information measure $I$ on the probabilities of events $p_i$ by the dependence of $\Lambda$ on the fundamental constants of nature (i.e., the fundamental coupling constants, masses and mixing angles), Eq. (16) was obtained from the requirements of ‘Fundamentality’, ‘Boundedness’, ‘Simplicity’ and ‘Scale-invariance’ [21].

A fourth derivation of Eq. (16) has been proposed in a recent series of papers by Burikham, Cheamsawat, Harko, Lake and Paterek [19, 20, 23, 24] (2016-2017). This is based on the construction of a dark energy-modified minimum length uncertainty relation, dubbed the dark energy uncertainty principle, or DE-UP for short, which may be combined with classical minimum mass, radius and/or density bounds to yield stability conditions for self-gravitating, compact and quantum mechanical objects. Applying the DE-UP to charge-neutral particles recovers the minimum-mass bound previously obtained in [16], which is consistent with current bounds on the mass of the electron neutrino [9], whereas applying it to charged particles with $Q = \pm e$ yields the electron mass, in accordance with Eq. (16).

The form of the DE-UP proposed in [19] leads naturally to holographic relation between the bulk and the de Sitter horizon, in which the number of minimum-volume ‘cells’ equals the number of Planck-sized ‘bits’ on the boundary [72, 73, 74]. However, it must be noted that, in order to maintain this relation for $\tau \ll \tau_0$, where $\tau_0 \simeq t_{dS} = l_{dS}/c$, we must substitute $l_{dS} \rightarrow r_H(\tau)$, where $r_H(\tau)$ is the physical horizon radius at cosmic time $\tau$. Performing the same substitution in the DE-UP – the validity of which is also supported by the physical arguments proposed in [20] – therefore gives rise to time-variation of the minimum mass for charge-neutral particles and of the maximum charge-squared to mass ratio for charged particles.
At present, further analysis is required to determine whether the predictions of the DE-UP, which gives rise to a natural implementation of the LNH at the present epoch, as well as to a natural description of late-time accelerated expansion in terms of dark energy particles (see Sec. 3.2), is compatible with existing cosmological data. However, it must be noted that the present model implicitly assumes that \( G, c, h \) and \( \Lambda \) are genuine universal constants which do not vary in time. Running of the gravitational coupling, as recently claimed in [37, 38, 39, 40], or additional time-dependence in any of the parameters \( \{G, c, h, \Lambda\} \), may dramatically alter these predictions. With this in mind, we note that a model in which \( \Lambda \propto \alpha_c^{-6} \), in the context of varying \( \alpha_c \) cosmology, was considered in [70, 71], whereas alternative ways of incorporating the effects of Universal expansion/dark energy on the uncertainty principle were considered in [75, 76].

Finally, we may consider the implications of a connection between “dark” and “light” physics, suggested by Eq. (16), for the physics of black holes. \( \textit{A priori} \), Eqs. (9)-(10) and (16) say nothing about black holes, yet if, as claimed in [19, 20], the ultimate origin of these relations is the DE-UP, it is by no means clear whether this even applies to objects with masses \( m \geq m_{Pl} \). In general, the form of positional uncertainty (if any) obeyed by the centre-of-mass of a black hole, remains an outstanding problem in contemporary theoretical physics. (See for [77, 78, 79, 80, 81, 82] for recent works in this direction.) Nonetheless, it is certainly worthwhile to attempt to extend the DE-UP into this region, which may be done naively by simply replacing the rest mass \( m \) with the ‘dual’ ADM mass \( m \rightarrow m_{ADM} \equiv m_{Pl}^2/m_{ADM} \approx m_{Pl}^2/(m + m_{Pl}^2/m) \). This gives rise to a unified Compton-Schwarzschild line connecting the black hole and particle regimes (see [77, 78, 79, 80, 81, 82] and [83, 84]). Since the DE-UP naturally implements holography in the \( m \leq m_{Pl} \) regime, it may be hoped that the extended version maintains it for \( m \geq m_{Pl} \), which may have profound implications for the black hole information loss paradox [85, 86, 87]. In this context, reassessing Beck’s information-theoretic approach, subject to holographic constraints, may prove particularly fruitful.

Acknowledgments
M.L. thanks the Institute for Fundamental Study (IF), Phitsanulok, Thailand, the Yukawa Institute for Theoretical Physics (YITP), Kyoto, Japan, and the Asia-Pacific Center for Theoretical Physics (APCTP), Pohang, Korea, for their generous support of the 6th IF+YITP International Symposium, held at Naresuan University in August 2016. He is supported by a Naresuan University Research Fund individual research grant.

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