Gravity lens critical test for gravity constants and dark sector

V. G. Gurzadyan\textsuperscript{1,2,a}, A. Stepanian\textsuperscript{1}

\textsuperscript{1} Center for Cosmology and Astrophysics, Alikhanian National Laboratory and Yerevan State University, Yerevan, Armenia
\textsuperscript{2} SIA, Sapienza Universita di Roma, Rome, Italy

Received: 14 August 2018 / Accepted: 18 October 2018 / Published online: 27 October 2018
© The Author(s) 2018

Abstract The recent study of the strong gravitational lens ESO 325-G004 (Collett et al., Science, 360:1342, 2018) leads to a new possibility for testing General Relativity and its extensions. Such gravity lens observational studies can be instrumental for establishing a limitation on the precision of testing General Relativity in the weak-field regime and on the two gravity constants (the Newtonian and cosmological ones) as described in Gurzadyan and Stepanian (Eur Phys J C 78:632 2018). Namely, we predict a critical value for the involved weak-field parameter $\gamma = 0.998$ (for $M = 1.5 \times 10^{11} M_\odot$ lens mass and $r = 2 kpc$ light impact distance), which remarkably does not depend on any hypothetical variable but is determined only by well measured quantities. If the critical parameter $\gamma_{cr}$ will be established at future observations, this will mark the first discrepancy with General Relativity of the conventional weak-field Newtonian limit, directly linked to the nature of the dark sector of the Universe.

1 Introduction

High precision tests of General Relativity (GR) and of its weak- and strong-field limits, along with of modified gravity models have gained new content with establishing of the dominating dark sector of the Universe. The models proposed to describe the dark energy and dark matter observations are being constrained by means of those tests. Weak-field General Relativity is continuously being tested in ever increasing accuracy \cite{3,4} including the frame-dragging measurements with laser ranging satellites \cite{5}.

The use of the strong gravity lens ESO 325-G004 \cite{1} demonstrates the efficiency of the lens studies to constrain the weak-field GR in the intergalactic scales. Given the increasing accuracy and statistics of the gravity lensing observations (e.g. \cite{6}), those studies can become laboratories to test GR modifications along with the structure and dynamics of the dark matter dominated configurations.

The idea of our approach to this problem is the following. In recent studies \cite{2,7} we have shown that the Newton theorem on the identity of the gravity of a sphere and of a point mass located in its center provides a natural way for the weak-field modification of GR. Consequently, the constant $\Lambda$ appears both in the cosmological solutions describing the accelerated expansion and at galactic halo scales cf. \cite{8}) as weak-field GR, thus linking dark energy and galactic scales.

Here we predict a critical value for the parameter of the parametrized post-Newtonian (PPN) formalism \cite{4} $\gamma = 0.998$ (normalized to given lens mass and light impact distance) which if observed at gravity lenses with proper significance will, for the first time, reveal the weak-field modification of GR differing from the conventional Newtonian limit. It is important that the critical $\gamma$ does not depend on any open parameter but involves only the fundamental constants and well measured quantities.

2 Newton’s theorem and gravity lensing with $\Lambda$

In \cite{2} it is shown that from the Newton theorem the weak-field limit of GR follows involving the cosmological constant $\Lambda$, so that the metric tensor components have the form

$$
g_{00} = 1 - \frac{2Gm}{rc^2} - \frac{\Lambda r^2}{3}; \quad g_{rr} = \left(1 - \frac{2Gm}{rc^2} - \frac{\Lambda r^2}{3}\right)^{-1}. \tag{1}
$$

While this metric (Schwarzschild–de Sitter) was obviously known before (e.g. \cite{10,11}), with the $\Lambda$ introduced by Einstein to get a static cosmological model, the motivation based on the Newton theorem essentially differs from that and ensures the Newton potential with the $\Lambda$ term as the weak-field limit of GR.
Namely, we take in our considerations the general function for the force satisfying Newton’s theorem on the identity of the sphere’s gravity and that of a point of the sphere’s mass situated in its center. That function besides the \( r^{-2} \) term contains also a second term \([9]\),

\[
f(r) = Ar^{-2} + \Lambda r. \tag{2}\]

When the modified Newtonian law (for the potential) is taken as the weak-field limit of GR, one has the constant \( \Lambda \) as the weak-field limit of GR, thus, the second constant acts as a cosmological constant in the solutions of Einstein equations, and at the same time it enters the low-energy limit of GR.

Turning to the strong lensing and following \([1]\), it is convenient to deal with the parameter representing the ratio

\[
\gamma = \Psi/\Phi \tag{3}
\]

of the two perturbing functions—of the Newtonian potential \( \Phi \) and the curvature potential \( \Psi \)—entering the weak-field metric

\[
ds^2 = (1 + 2\Phi)c^2 dt^2 - (1 - 2\Psi)r^2 dr^2 - r^2 d\Omega^2. \tag{4}\]

For GR \( \gamma = 1 \), obviously.

In the weak-field limit of GR following from the Newton theorem Eq. \((1)\) \([2,7]\), we obtain

\[
\Phi = \Psi = \frac{GM}{rc^2} - \frac{\Lambda r^2}{6}. \tag{5}\]

To find \( \gamma \) the authors of \([1]\) introduced the following dynamical mass of the lens:

\[
M_{\text{dyn}} = \frac{1 + \gamma}{2} M, \tag{6}\]

where consequently the light bending angle in the vicinity of a mass of surface density \( \Sigma \) is obtained:

\[
\alpha = \frac{2(1 + \gamma)}{c^2} \int d^2 x' \Sigma(x') \frac{x - x'}{|x - x'|^2}. \tag{7}\]

Upon the analysis of the observational data of ESO 325-G004 they finally obtain the value \( \gamma \simeq 0.97 \pm 0.09 \).

Now, within our approach of the weak-field metric of Eq. \((1)\) the bending angle will be \([12]\]

\[
\alpha = \frac{4GM}{c^2 r} - \frac{\Lambda c^2 r^3}{6GM}. \tag{8}\]

Here, an important difference arises between the Newtonian and the \( \Lambda \)-modified cases. Namely, the authors in \([1]\) have obtained

\[
\alpha = 2(1 + \gamma) \frac{GM}{c^2 r}. \tag{9}\]

Comparing this with Eq. \((8)\), we get for the \( \gamma \)-parameter

\[
\gamma = 1 - \frac{\Lambda c^4 r^4}{12G^2 M^2}. \tag{10}\]

Inserting the current value of the cosmological constant e.g. that of the Planck satellite \([13]\) \( \Lambda = 1.11 \times 10^{-52} \text{ m}^{-2} \) we obtain

\[
\gamma_{cr} = 1 - 0.002 = 0.998 \left( \frac{M}{1.5 \times 10^{11} \text{ M}_\odot} \right)^{-2} \left( \frac{r}{2 \text{ kpc}} \right)^4, \tag{11}\]

where the data of \([1,14]\) for ESO 325-G004, i.e. the Einstein radius and the estimated mass inside that radius, were used for the normalization. Obviously, the normalization and hence the precise numerical value of \( \gamma_{cr} \) will vary from one lens to another, and the principal point is the existence of a well defined parameter enabling one to reveal the weak-field limit of GR with the lensing effect.

The other key point of using ESO 325-G004 is that within other available observational means it is technically impossible to detect the contribution of the \( \Lambda \) term in the gravitational lensing. For example, for the same effect within the Solar System the value of \( \gamma_{cr} \) will be approximately \((1-9.6 \times 10^{-25})\). Note that a limitation on \( \gamma \) also will emerge due to the proper motion of the lens, for it affects the measured value of \( \gamma \) as shown in \([15]\).

### 3 Conclusions

The accurate measurements of strong lensing of extragalactic objects provide important means for the study of profound cosmological problems. The recent study of the lensed object ESO 325-G004 \([1]\) (cf. previous studies \([6,16,17]\)) enables a remarkable test of General Relativity on extragalactic scales. The observational surveys of lensing will definitely proceed with ever increasing precision and statistics, which will enable one to improve the available accuracy of the value of the weak-field parameter \( \gamma \).

In view of that, here we derive a critical \( \gamma_{cr} \) for the strong lensing of extragalactic objects which can be informative for gravity theories. It is remarkable that \( \gamma_{cr} \) does not depend on any hypothetical parameter of modified gravity models (coupling constant, scalar field mass, etc.) but is determined entirely by measured physical quantities. Since the needed accuracy for measuring of \( \gamma_{cr} \) does not seem to be principally unreachable given the variability range of the lens mass and light impact scale, due to \( \gamma_{cr} \) the gravity lens measurements
will get a similar importance to the renowned Solar eclipse of 1919, which enabled one to distinguish GR from classical Newtonian gravity. Thus, for the first time the detection of a discrepancy with the conventional General Relativity can become feasible, with a further intriguing relation to the dark sector. While the breakthrough study [1] reveals the possibility of obtaining $\gamma$, the observations of more distant objects and hence the needed accuracy for $\gamma$ certainly is a matter of future advances; however, let us recall the classical example, i.e. Einstein’s skepticism as regards observing gravitational lenses [18].

Acknowledgements We are thankful to the referee for helpful comments and to S. Mirzoyan for discussions. AS acknowledges the Abdus Salam International Centre for Theoretical Physics Affiliated Center program AF-04 for financial support.

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. Funded by SCOAP3.

References

1. T.E. Collett et al., Science 360, 1342 (2018)
2. V.G. Gurzadyan, A. Stepanyan, Eur Phys. J. C 78, 632 (2018)
3. S. Kopeikin, M. Efroimsky, G. Kaplan, Relativistic celestial mechanics of the solar system (Wiley, New York, 2001)
4. C.M. Will, Living Rev. Relativ. 17, 4 (2014)
5. I. Ciufolini et al., Eur. Phys. J. C 76, 120 (2016)
6. S. Cao, APJ 835, 92 (2017)
7. V.G. Gurzadyan, arXiv:1712.10014 (2017)
8. V.G. Gurzadyan et al., A & A 609, A131 (2018)
9. V.G. Gurzadyan, Observatory 105, 42 (1985)
10. W. Rindler, Relativity, special, general, and cosmological (Oxford University Press, Oxford, 2006)
11. M. Nowakowski, Int. J. Mod. Phys. D 10, 649 (2001)
12. M. Ishak, W. Rindler, J. Dossett, J. Moldenhauer, C. Allison, MNRAS 388, 1279 (2008)
13. P.A.R. Ade et al., A&A 594, A13 (2016)
14. R.J. Smith, J.R. Lucey, MNRAS 434, 1964 (2013)
15. S. Kopeikin, MNRAS 399, 1539 (2009)
16. A.S. Bolton, S.A. Rappaport, S. Burles, Phys. Rev. D 74, 061501 (2006)
17. J. Schwab, A.S. Bolton, S.A. Rappaport, APJ 708, 750 (2010)
18. A. Einstein, Science 84, 506 (1936)