Modelling the X-ray polarimetric signatures of complex geometry: the case study of the “changing look” AGN NGC 1365

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ABSTRACT

“Changing look” Active Galactic Nuclei (AGN) are a subset of Seyfert galaxies characterized by rapid transitions between Compton-thin and Compton-thick regimes. In their Compton-thin state, the central engine is less obscured, hence spectroscopy or timing observations can probe their innermost structures. However, it is not clear if the observed emission features and the Compton hump are associated with relativistic reflection onto the accretion disc, or complex absorption by distant, absorbing gas clouds passing by the observer’s line-of-sight. Here, we investigate these two scenarios under the scope of X-ray polarimetry, providing the first polarisation predictions for an archetypal “changing look” AGN: NGC 1365. We explore the resulting polarisation emerging from lamp-post emission and scattering off an accretion disc in the immediate vicinity of a supermassive black hole. The computed polarisation signatures are compared to the results of an absorption-dominated model, where high column density gas partially covers the central source. While the shape of the polarisation spectrum is similar, the two models differ in net polarisation percentage, with the relativistic reflection scenario producing significantly stronger polarisation. Additionally, the variation of the polarisation position angle is distinctly different between both scenarios: the reflection-dominated model produces smooth rotations of the polarisation angle with photon energy whereas circumnuclear absorption causes an orthogonal switch of the polarisation angle between the soft and the hard X-ray bands. By comparing the predicted polarisation of NGC 1365 to the detectability levels of X-ray polarimetry mission concepts proposed in the past, we demonstrate that with a large, soft X-ray observatory or a medium-sized mission equipped with a hard (6 – 35 keV) polarimeter, the correct interpretation would be unambiguous.

Key words: polarisation – radiative transfer – line: profiles – scattering – X-rays: galaxies – galaxies: active.

1 INTRODUCTION

AGN are complex objects, hosting several media over a broad range of location. Radiation from the central, ionising source is absorbed and reprocessed by a variety of scattering regions such as the accretion disc, the Broad Line Region (BLR), the Warm Absorber (WA), the molecular torus, and the Narrow Line Region (NLR), imprinting the resulting X-ray spectrum of AGN. For example, the type 1.8 AGN NGC 1365 (z = 0.00547), also known as the Great Barred Spiral Galaxy, is an intriguing source displaying the most dramatic X-ray spectral changes observed so far in an AGN: the source switched from reflection-dominated (Compton-thick) to transmission-dominated (Compton-thin) and back in just a few days to week time-scales (Risaliti et al. 2005). NGC 1365 can be classified as one of the most extreme examples of “changing look” AGN. This has been attributed to variation in the line-of-sight of cold absorber rather than to extreme intrinsic emission variability. Such rapid timescales of these X-ray eclipses suggest that these absorbers
(N$_H \sim 10^{23} – 10^{24}$ cm$^{-2}$) are located on compact scales consistent within the inner BLR and the outer part of the accretion disc. Such behaviour is now observed in numerous other sources, e.g. NGC 4388 (Elvis et al. 2004), NGC 7674 (Bianchi et al. 2005), NGC 4151 (Puccetti et al. 2007), NGC 7582 (Bianchi et al. 2009), UGC 4203 (Guainazzi et al. 2002; Risaliti et al. 2010), NGC 4051 (Guainazzi et al. 1998; Lobban et al. 2011), 1H 0419-577 (Pounds et al. 2004), NGC 454 (Marchese et al. 2012), IRAS 09104+4109 (Piconcelli et al. 2007; Chiang et al. 2013), H0557-385 (Longinotti et al. 2009), and the Phoenix galaxy (Matt et al. 2009). Besides, it has been shown that NGC 1365 exhibits an apparent broad Fe K$_\alpha$ fluorescent emission line (Brenneman et al. 2013; Risaliti et al. 2013). Two alternative origins are proposed to explain the shape and variability of this feature, general relativistic effects (Risaliti et al. 2013) or partial-covering absorptions (Miller & Turner 2013). In the former scenario, reflection of hard X-rays off the inner disc, reaching down to the innermost stable circular orbit (isco), produces fluorescent emission, subsequently broadened by general relativistic and Doppler effects. Integrating the reprocessed emission across the disc constitutes the blurred line that is sensitive to the spin of the SMBH (Fabian et al. 1989, Laor 1991, Dovčiak et al. 2004, Brenneman & Reynolds 2006, Dauser et al. 2010). In the latter scenario, complex absorption due to clouds of absorbing gas, partially covering the central AGN source, is responsible for both the flux variability and the apparent broadening of the Fe line (Inoue & Matsumoto 2003, Miller, Turner & Reeves 2008, 2009). Continuum radiation transmitted and scattered through circumnuclear medium carves out the extended red wing, which leads to systematically lower measurements of the SMBH spin than obtained by the relativistic reflection method. Interestingly, the variability behaviour of the Fe K line is contrary to that of the relativistic line seen in the type-1 Seyfert galaxy MCG-6-30-15 (Fabian et al. 2002), for which light bending effect has been proposed (Fabian & Vaughan 2003; Miniutti et al. 2003). However, this can alternatively be explained by an increase of the column density of the cold absorber from 2008 to 2010 that would reduce the red wing of the line as pinpointed by Brenneman et al. (2013).

In this context, we aim to illustrate the potential of X-ray polarimetry observations as a powerful tool to test the complex geometry in “Changing look” AGN, using NGC 1365 as a study case.

2 POLARISATION INDUCED BY RELATIVISTIC REFLECTION

The combined XMM-Newton / NuSTAR observation of NGC 1365 (Risaliti et al. 2013) puts forward a scenario in which the Fe K$_\alpha$ emission line originates from a cold accretion disc surrounding a SMBH of mass $M = 2 \times 10^6$ M$_\odot$, associated with a dimensionless spin parameter $\alpha = 1$ (extreme Kerr BH). Their model assumes that the X-ray source, illuminating the accretion disc, is localised on the rotation axis, at 2.5 gravitational radii above the black hole. The broadening of the Fe K$_\alpha$ fluorescence line is then a consequence of rapid motion and general relativistic effects close to the SMBH. Using the method already described in details in Dovčiak et al. (2011) and used in Marin et al. (2012a) to model X-ray polarisation spectra emerging from the Seyfert-1 MCG-6-30-15, we now apply it to the case of “changing look” AGN in their Compton-thin regime.

Relying on the prescribed reflection model (Risaliti et al. 2013), we run Monte Carlo radiative transfer simulations of the $1 – 100$ keV polarisation emerging from NGC 1365. The input spectrum ($> 10^9$ photons) has a power-law shape $F_\nu \propto \nu^{-\alpha}$ with a spectral index $\alpha = 1.2$ (Risaliti et al. 2009, Risaliti et al. 2013). The intrinsic flux is unpolarised and emitted isotropically. The re-emitted intensity as a function of incident and re-emission angle is computed by the Monte-Carlo radiative transfer code NOAR (Dumont et al. 2000), the local polarisation being estimated according to the transfer equations of Chandrasekhar (1960). The local, polarised reflection spectra are then combined with the KY-code (Dovčiak et al. 2004), that conducts relativistic ray-tracing between the elevated source, the disc, and the distant observer (see Fig. 1, left-hand panel).

The resulting polarisation signatures, as a function of photon energy, are plotted in Fig. 2 (red dashed line). $P$ is the degree of linear polarisation emerging from the models and $\Delta \Psi$ the rotation of the polarisation position angle with respect to a convenient average of the polarisation position angle over the depicted energy band. An “absolute orientation” of the polarisation plane is difficult to measure, and we thus focus on the variation of $\Delta \Psi$ with photon energy.

Considering an inclination of $i = 60^\circ$ for NGC 1365 (Risaliti et al. 2013), the reflection model is found to produce relatively high degrees of polarisation ($P \geq 1 \%$) in the $3 – 100$ keV band, with a maximum $P$ in the Compton hump, where multiple scattering dominates. Below 10 keV, $P$ decreases sharply, unlike an energy-independent Newtonian case, as the local disc polarisation is affected by the relativistic energy shift of the photons and impacts the resulting polarisation seen by a distant observer (Dovčiak et al. 2004). The energy-dependence of $P$ in the soft X-ray band is accentuated by a dilution mechanism caused by the source, set to favour the production of soft X-ray photons: even if strong gravity effects tend to bend photons back to the disc, a fraction of the input, unpolarised radiation merges with the (reprocessed) polarised photons and dilutes the $P$ signal. Thus, the minimum polarisation degree produced by disc reflection is $\geq 0.1 \%$ above 1 keV, $\geq 1 \%$ between 3 and 10 keV, and up to $10 \%$ in the Compton hump regime.

Interestingly (Fig. 2, bottom, red dashed line), the relativistic reflection scenario exhibits an energy-dependent $\Delta \Psi$ that varies significantly across the Fe K$_\alpha$ line, a behaviour already expected from X-ray polarisation simulations in the type-1 AGN MCG-6-30-15 (Marin et al. 2012a). The feature is related to the energy-dependent albedo and scattering phase function of the disc material. $\Delta \Psi$ presents a smooth and continuous variation by $\sim 5^\circ$ around the iron line and increases monotonically upward 20 keV.

3 PURE ABSORPTION AND COMPTON SCATTERING BY A CLOUDY MEDIUM

Taking into account that the changing nature of NGC 1365 is explained by the presence of a clumpy distribution of gas around the central engine, we now investigate the result-
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The central, X-ray source is defined as a geometrically thin, emitting slab that may represent the so-called hot inner flow. The emitting region is insensitive to scattering and does not reach down to the ISCO in order to avoid a possible confusion with relativistic disc reflection. Inside a spherical shell constrained between $r_1 = 10$ and $r_2 = 20$ (in units of spheres radius) from the model’s origin, a random distribution of 1 000 scale-independent, Compton-thick, absorbing spheres with equal radius and constant density, was generated. The covering factor of the cloud distribution varies systematically from the poles towards the equator, as suggested by the unified scheme of AGN, and is $\sim 0.55$ for $i = 0^\circ$, and $> 0.75$ for $i > 60^\circ$ (Antonucci 1993; Urry & Padovani 1995). The column density of the absorbing, cold medium in a single cloud is equal to $1.5 \times 10^{24}$ cm$^{-2}$, a value chosen to maximise the observed hard excess. For coherency, we restrain ourselves to the same input power-law spectrum as in Sect. 2. This parametrisation is in perfect agreement with the modelling of Miller & Turner (2013), allowing us, in the mean time, to examine the net polarisation of “changing look” AGN and explore the polarimetric signatures of their theoretical, absorption-dominated model.

We find that, in the case of an absorption-dominated model (Fig. 2 top), the resulting polarisation degree is energy-dependent (solid black line), with a maximum $P$ observed in the Compton hump ($P \lesssim 0.5 \%$). Close to 15 keV, the local polarisation strongly diminishes due to a flip of the photon position angle, as perpendicular and parallel photon’s position angles cancel each other (see next paragraph). Finally, $P$ gradually decreases below the iron fluorescence line energy, as the dilution trend is strengthened by the absence of strong gravity effects bending photons back to the disc. As forward scattering dominates, we naturally find that the polarisation produced by a clumpy distribution of cold matter is about 15 times weaker than the relativistic model.

The most interesting results come from the variation of the polarisation position angle (Fig. 2 bottom). The cloud distribution examined in this Letter (solid black line) shows...
an instantaneous $\Delta \Psi$ transition by $90^\circ$ from the soft (below 15 keV) to the hard X-ray band, associated with a local diminution of $P$. In the 1 – 15 keV band, radiation is not energetic enough to scatter through the highly covered equatorial plane and becomes absorbed. Photons escaping from the system scatter on polar clouds, with their normal to the scattering plane parallel to the axis of the system. Thus, the polarisation vectors in the equatorial plane of the system, aligning their electric field vector to the projected axis of the system. Soft X-rays (red lines) are absorbed along the highly covered equatorial plane; to escape, radiation reprocesses on polar clouds and the average direction of the electric field vector of the scattered radiation is perpendicular to the axis of the system. Hard X-rays (green line) face a lower electron scattering cross-section and easily escape the circumnuclear medium by multiple scattering, aligning their electric field vector to the projected axis of the system. Thus, a $\Psi$ rotation between the two regimes occurs at a model-dependent energy.

4 OBSERVATIONAL PROSPECTS

We showed, in a previous publication centred around MCG-6-30-15 (Marin et al. 2012b), that a small pathfinder mission could be able to distinguish between relativistic reflection and transmission through absorbing media. But, is this still true for NGC 1365, a rather different object with a much higher ($60^\circ$) inclination, a variable behaviour of its iron line red-wing, and extreme Compton-thick to Compton-thin regime transitions? Taking into account an average flux in the X-ray band ($\sim$ 1.0 mCrab in the 2 – 10 keV band, Brenneman et al. 2013, and $\sim$ 2.62 ± 0.50 mCrab in the 17 – 60 keV band, Krivonos et al. 2007), we extend the sample of already proposed X-ray mission concepts of different sizes which included a polarimeter to investigate if any polarisation signal could be detected.

- **XIPE**, the S-class X-ray Imaging Polarimetry Explorer (Soffitta et al. accepted), designed to be loaded with two 2 – 10 keV Gas Pixel Detectors (GPD, Costa et al. 2001, Bellazzini et al. 2006, Bellazzini & Muleri 2010).
- The M-class New Hard X-ray Mission **NHXM** (Tagliaferri et al. 2012a,b). One of the four telescopes aboard would be coupled with two GPD instruments: a low energy polarimeter (LEP, 2 –10 keV) and a medium energy polarimeter (MEP, 6 – 35 keV), for a total coverage of 2 – 35 keV.
- The L-class International X-ray Observatory **IXO** (Barcons et al. 2011) included, among a few other instruments, the 2 – 10 keV X-ray polarimeter XPOL (Bellazzini et al. 2006).

In Fig. 4 we calculated the representative minimum detectable polarisation (MDP, whose formula can be found in

\begin{equation}
  \text{MDP} = \frac{2 \times \text{Sensitivity}}{\text{Flux}}
\end{equation}

Figure 3. Schematic illustration of the dominant polarisation mechanism for absorption and Compton scattering by a cloudy medium. Soft X-rays (red lines) are absorbed along the highly covered equatorial plane; to escape, radiation reprocesses on polar clouds and the average direction of the electric field vector of the scattered radiation is perpendicular to the axis of the system. Hard X-rays (green line) face a lower electron scattering cross-section and easily escape the circumnuclear medium by multiple scattering, aligning their electric field vector to the projected axis of the system. Thus, a $\Psi$ rotation between the two regimes occurs at a model-dependent energy.

Figure 4. Representative minimum detectable polarisation (MDP) of the two scenarios for a 1 Ms observation of NGC 1365. The observer’s line-of-sight lies at $60^\circ$ with respect to the symmetry axis. The MDP for **XIPE** is represented by the maroon line, for **NHXM** in blue and for **IXO** in green. Legend: a fragmented absorption region (solid line) and a relativistic reflection model with an extreme Kerr SMBH with $a = 1$ (red dashed line).
Marin et al. 2012a) on the whole energy range of the instrument at 99 % confidence level that the mission concepts would have reached pointing at NGC 1365. The observation time is estimated to 1 Ms, taking into account that the background flux is negligible with respect to the source flux. We assumed a Crab-like spectrum, that is a power law with an index of 2.05 and fluxes reported above. The rate of the source is calculated on a fine energy grid (0.1 keV) by multiplying the telescope collecting area, the detector efficiency and the source spectrum. Then, we multiply (for each energy) the square root of the rate by the modulation factor, adding the background (which is negligible by orders of magnitude in this case) and smearing the product for the energy resolution of the instrument. Eventually, we sum up the values of the product in the energy range of interest to derive the MDP for the selected interval. The MDP evaluated over representative energy bands are: $\sim 7 \%$ ($4 \text{ – } 10 \text{ keV}$) for XIME, $\sim 3 \%$ for MEP/NHXM ($6 \text{ – } 35 \text{ keV}$) and $\sim 0.3 \%$ for IXO ($2 \text{ – } 10 \text{ keV}$).

We find that the predicted polarisation signal induced by relativistic reflection is below the polarisation detectability of XIME (maroon line) so, for an exposure time of 1 Ms, a small, pathfinder mission such as XIME could not measure the polarisation unless the observed polarisation is significantly (a factor of 2) higher than that predicted by the model. However, any detection would strongly support the relativistic model that produces a conservative polarisation degree (see Sect. 3). Further indications could be deduced from the variation of the polarisation position angle around the iron line if the observed polarisation signal is significant enough. The blue line, representing the MDP of MEP/NHXM, indicates that the polarisation signal originating from the relativistic reflection case can be detected across the whole $6 \text{ – } 35 \text{ keV}$ band. It covers in particular the iron line domain, where the $\Delta \Psi$ signature of strong gravity would be detected. Finally, for longer observing time or higher flux levels, IXO (green line) would have been the only mission to potentially detect the low polarisation originating from a “changing look” AGN in the complex absorption scenario.

5 DISCUSSION AND CONCLUSIONS

Due to its transient nature, the distribution of gas clouds around the central region of NGC 1365 may vary, both in its covering factor and column density. Lowering the covering factor would result in a decrement of the net polarisation as dilution from the unscattered input spectrum would be enhanced. A different density parameterisation would also affect the polarisation degree: optically thinner reprocessing clouds would reduce the probabilities of scattering, thus decreasing $P$, while increasing the gas density would slightly increase the polarised flux until the circumnuclear matter becomes optically too thick for the radiation to escape. Nonetheless, since absorption-dominated models are based upon continuum transmission through gas, the overall polarisation degree is expected to be weaker than for (perpendicular) scattering off an accretion disc. One must note that the energy at which the $\Psi$ rotation occurs is also governed by the gas density, covering factor, and possibly inclination. We intend to investigate different realisations of a variety of “changing look” AGN in future work in order to explore the complex inter-combination of parameters that may lead to similar $\Psi$ rotation. This feature, if observable, could be used as a major constraint to probe the parameter space of the cold gas.

The relativistic reflection model employed to simulate the polarisation response of NGC 1365 is basic, yet instructive. In a forthcoming publication, we intend to improve our relativistic models by taking into account radially structured surface ionisation and intrinsically polarised, non-axisymmetric irradiation. Both implementations are expected to increase the net polarisation, similarly to coronal emission that scatters off an ionised disc and disc return radiation over the potential well (Schnittman & Krolik 2010; 2013). Therefore, the modelling presented in this Letter produces a conservative, lower limit on the polarisation degree.

The polarisation position angle across the iron line turns out to be a particularly strong indicator for the correct scenario – if it can be constrained with a sufficiently powerful X-ray polarimetry mission and/or long enough observations. The bottom panel of Fig. 2 shows that in the case of relativistic reflection $\Delta \Psi$ follows a characteristic “S-shape” between 2 and 9 keV, whereas in the absorption scenario the distribution of $\Delta \Psi$ is only bipolar. The detection of a few degrees swing could be within the reach of a large mission, for which the MDP is significantly lower than the expected degree of polarisation. The fact that in the absorption case only “parallel” or “perpendicular” polarisation are allowed is related to the axial symmetry of the cloud distribution. The switch exists due to a change in the scattering albedo between the soft and the hard bands of the X-ray spectrum. Goosmann & Matt (2011) showed that a more gradual rotation of $\Delta \Psi$ can be obtained in a non-relativistic scenario once the distribution of scattering clouds is no longer symmetric. However, in this case the rotation can only evolve monotonically with photon energy between two limiting values determined by the net soft and hard X-ray polarisation states. Such a scenario could not reproduce the S-shape variation of $\Delta \Psi$ seen in the relativistic model.

In conclusion, this work predicts the polarisation degree expected from “changing look” AGN in their Compton-thin state. The reflection-dominated scenario of NGC 1365 is found to be associated with strong polarisation degrees and a smooth variation of polarisation angle across the iron line emission. However, if the spectrum of NGC 1365 is dominated by cloud absorption, lower polarisation degrees are expected, associated with an orthogonal switch of the polarisation position angle between the soft and the hard X-ray band. We strengthen the conclusions drawn in Marin et al. (2012a), despite some fundamental differences between MCG-6-30-15 and NGC 1365, and demonstrate that X-ray polarimetry measurement of NGC 1365 could discriminate between the two scenario using a hard polarimeter or a large, soft X-ray polarimetric mission.

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REFERENCES

Antonucci, R. 1993, ARA&A, 31, 473
Barcons, X., Barret, D., Bautz, M., et al. 2011, arXiv:1102.2845
Bellazzini, R., & Muleri, F. 2010, Nuclear Instruments and Methods in Physics Research A, 623, 766
Bellazzini, R., Baldini, L., Brez, A., et al. 2006, Proceedings of the SPIE, 6206
Bianchi, S., Guainazzi, M., Matt, G., et al. 2005, A&A, 442, 185
Bianchi, S., Piconcelli, E., Chiaberge, M., et al. 2009, ApJ, 695, 781
Brenneman, L. W., & Reynolds, C. S. 2006, ApJ, 652, 1028
Brenneman, L. W., Risaliti, G., Elvis, M., & Nardini, E. 2013, MNRAS, 429, 2662
Chandrasekhar, S. 1960, New York: Dover, 1960
Chiang, C.-Y., Cackett, E. M., Gandhi, P., & Fabian, A. C. 2013, MNRAS, 430, 2943
Costa, E., Soffitta, P., Bellazzini, R., et al. 2001, Nature, 411, 662
Dauser, T., Garcia, J., Wilms, J., et al. 2013, MNRAS, 430, 1694
Dovciak, M., Karas, V., & Yaqoob, T. 2004, ApJS, 153, 205
Dovciak, M., Muleri, F., Goosmann, R. W., Karas, V., & Matt, G. 2011, ApJ, 731, 75
Dumont, A.-M., Abrassart, A., & Collin, S. 2000, AAP, 357, 823
Elvis, M., Risaliti, G., Nicastro, F., et al. 2004, ApJL, 615, L25
Fabian, A. C., Rees, M. J., Stella, L., & White, N. E. 1989, MNRAS, 238, 729
Fabian, A. C., Vaughan, S., Nandra, K., et al. 2002, MNRAS, 335, L1
Fabian, A. C., & Vaughan, S. 2003, MNRAS, 340, L28
Goosmann, R. W., & Matt, G. 2011, MNRAS, 415, 3119
Goosmann, R. W., & Gaskell, C. M. 2007, A&A, 465, 129
Guainazzi, M., Nicastro, F., Fiore, F., et al. 1998, MNRAS, 301, L1
Guainazzi, M., Matt, G., Fiore, F., & Perola, G. C. 2002, A&A, 388, 787
Inoue, H., & Matsumoto, C. 2003, PASJ, 55, 625
Karas, V., Czerny, B., Abrassart, A., & Abramowicz, M. A. 2000, MNRAS, 318, 547
Krivonos, R., Revnivtsev, M., Lutovinov, A., et al. 2007, A&A, 475, 775
Laor, A. 1991, ApJ, 376, 90
Lobban, A. P., Reeves, J. N., Miller, L., et al. 2011, MNRAS, 414, 1965
Longinotti, A. L., Bianchi, S., Ballo, L., de La Calle, I., & Guainazzi, M. 2009, MNRAS, 394, L1
Marchese, E., Braito, V., Della Ceca, R., Caccianiga, A., & Severgnini, P. 2012, MNRAS, 421, 1803
Marin, F., & Goosmann, R. W. 2011, SF2A-2011: Proceedings of the Annual meeting of the French Society of Astronomy and Astrophysics, 597
Marin, F., Goosmann, R. W., Dovcik, M., et al. 2012a, MNRAS, 426, L101
Marin, F., Goosmann, R. W., Gaskell, C. M., Porquet, D., & Dovcik, M. 2012b A&A, 548, A121
Matt, G., Bianchi, S., Awaki, H., et al. 2009, A&A, 496, 653
Miller, L., Turner, T. J., & Reeves, J. N. 2008, A&A, 483, 437
Miller, L., Turner, T. J., & Reeves, J. N. 2009, MNRAS, 399, L69
Miller, L., & Turner, T. J. 2013, ApJL, 773, L5
Minutti, G., Fabian, A. C., Goyder, R., & Lasenby, A. N. 2003, MNRAS, 344, L22
Piconcelli, E., Fiore, F., Nicastro, F., et al. 2007, A&A, 473, 85
Pounds, K. A., Reeves, J. N., Page, K. L., & O’Brien, P. T. 2004, ApJ, 605, 670
Puccetti, S., Fiore, F., Risaliti, G., et al. 2007, MNRAS, 377, 607
Risaliti, G., Elvis, M., Fabbiano, G., Baldi, A., & Zezas, A. 2005, ApJL, 623, L93
Risaliti, G., Braito, V., Laparola, V., et al. 2009, ApJL, 705, L1
Risaliti, G., Elvis, M., Bianchi, S., & Matt, G. 2010, MNRAS, 406, L20
Risaliti, G., Harrison, F. A., Madsen, K. K., et al. 2013, Nature, 494, 449
Schnittman, J. D., & Krolik, J. H. 2010, ApJ, 712, 908
Schnittman, J. D., & Krolik, J. H. 2013, arXiv:1302.3214
Soffitta, P., et al. 2013, Experimental Astronomy, in press
Tagliaferri, G., Hornstrup, A., Huovelin, J., et al. 2012a, Experimental Astronomy, 34, 463
Tagliaferri, G., & NHXM consortium, o. b. o. t. 2012b, MemSAI, 83, 360
Urry, C. M., & Padovani, P. 1995, PASP, 107, 803

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