The X-ray emission from Nova V382 Velorum: II. The super-soft component observed with BeppoSAX

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\textbf{ABSTRACT}

Nova Velorum 1999 (V382 Vel) was observed by BeppoSAX 6 months after optical maximum and was detected as a bright X-ray supersoft source, with a count rate 3.454±0.002 cts s\(^{-1}\) in the LECS. It was the softest and most luminous supersoft source observed with this instrument. The flux in the 0.1–0.7 keV range was not constant during the observation. It dropped by a factor of 2 in less than 1.5 hour and then was faint for at least 15 minutes, without significant spectral changes. The observed spectrum is not well fit with atmospheric models of a hot, hydrogen burning white dwarf. This is due mainly to a supersoft excess in the range 0.1–0.2 keV, but the fit can be significantly improved at higher energy if at least one emission feature is superimposed. We suggest that a “pseudocontinuum” was detected, consisting of emission lines in the supersoft X-ray range superimposed on the thermal continuum of a white dwarf atmosphere. As a result, an accurate determination of the effective temperature and gravity of the white dwarf at this post-outburst stage is not possible.

\textbf{Key words:} Stars: individual: V382 Vel, novae, cataclysmic variables – Sources as function of wavelength: X-rays: stars

1 INTRODUCTION

Nova Velorum 1999 (V382 Vel) was the second brightest nova of the last 50 years (V≃2.6) (Seargent & Pearce, 1999). It was a “fast ” O–Ne–Mg nova, with \(v_{\text{nova}}\) of the last 50 years (V=2.6) (Seargent & Pearce, 1999). Nova Velorum 1999 (V382 Vel) was the second brightest nova of the last 50 years (V=2.6) (Seargent & Pearce, 1999).

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Nova Velorum 1999 (V382 Vel) was observed with BeppoSAX, with ASCA and RossiXTE two weeks after the outburst (Orio et al. 1990, Orio et al. 2001b, Mukai & Ishida 1999, 2001) as a hard X-ray source (with plasma temperature \(kT\lesssim7\) keV). It cooled rapidly in the first two months after outburst to \(kT\lesssim2.4\) keV (Mukai & Ishida 2001). In a recent paper (Orio et al. 2001b; hereafter, Paper I) we attributed the initial X-ray emission to shocks in the nebula. Since the initially very large intrinsic absorption of the ejected nebula was thinning out, the equivalent N(H) decreased rapidly (Mukai & Ishida 2001). Thus, in the second BeppoSAX observation we hoped to detect the super-soft X-ray emission with the LECS and derive useful information on the nature of the white dwarf.

After a few months, luminous “supersoft” X-ray emission (luminosity of the order \(10^{35–36}\) erg s\(^{-1}\)) has also been observed. The previous detections have been attributed to residual hydrogen burning in a shell on the white dwarf remnant (e.g. Ögelman et al. 1993, Krautter et al. 1996, Orio & Greiner 1999). We expect to detect in this case an atmospheric continuum at \(T_{\text{eff}}\approx20-80\) eV and the absorption edges of the white dwarf (or even emission edges if the effective temperature is extremely high).

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The peak temperature of the hot white dwarf remnant, and the absorption edges that indicate the underlying chemical composition, are extremely important in order to constrain the physical models. In addition to this, the length of the constant bolometric luminosity phase is an essential parameter in order to understand whether the nova retains accreted mass after the outburst and is therefore a candidate type Ia supernova (or, even, a candidate neutron star formed by accretion induced collapse).

2 THE SUPERSOFT X-RAY SOURCE

In the second BeppoSAX observation on 1999 November 23, the actual length of the LECS exposure was 12.4 ksec, over about a total 16 hours in 11 time intervals, lasting for t≤1200 sec each. The MECS exposure time was 25.9 ksec during the same 16 hours. The count rate measured with the LECS was extremely high, 3.45±0.0021 cts s⁻¹ in the 0.1-4.0 keV range, due to the emergence of supersoft X-ray emission (Orio et al. 1999b). This high count rate was not unexpected, and it would be equivalent (assuming for simplicity a blackbody at 40 eV and N(H)=2×10¹² cm⁻²) to ≈50 cts s⁻¹ with the ROSAT PSPC. A count rate of 75 cts s⁻¹ was measured for V1974 Cyg at maximum with the PSPC (Krautter et al. 1996). In the 0.8-10.0 keV range the count rate was only 0.103±0.0038, consistent with the MECS count rate 0.045±0.0015 cts s⁻¹ (more than a factor of 3 lower than in June of 1999). There was no PDS detection with a 2σ upper limit of 0.080 cts s⁻¹ in the 15-50 keV range. We have already discussed the evolution of the hard X-ray emission (Paper I). We found that more than one component was necessary to fit the LECS spectrum, however we also concluded that there was no component with a higher plasma temperature than ≈1 keV, that N(H) was consistent with the interstellar value and that the supersoft portion of the flux was dominant.

Remarkably, the supersoft X-ray flux (0.1-0.7 keV) was variable. No significant variability was detected at higher energy. Overall, there was irregular flickering with time scales of minutes, and as Fig. 1 shows, in the 9th observation (after ≈13 hours from the beginning of the exposures) the background-subtracted count rate decreased dramatically, by approximately a factor of 2. This low state lasted during the whole LECS coverage of 15 minutes, spaced about 5000 seconds from two observation in which the average count rate was twice higher. Close to the end of the ≈16 hours of intermittent observations, the count rate decreased again in the few minutes (see bottom panel of Fig. 1). If we missed other episodes of this type, they must have been shorter than the 4000-5000 seconds that elapsed between the observations. One possible explanation for the sudden decrease in count rate is of course that a dense clump intervened along the line of sight. However, the spectrum was definitely supersoft during the whole observation and became slightly softer during the dip. Instead, the additional absorption of a thick clump would absorb the softer portion of the spectrum more and produce an apparently “harder” spectrum.

Is this phenomenon linked with orbital variability? Given the observed periodicity at optical wavelengths (Bos et al. 2001) the orbital period of V382 Vel is likely to be 3.5 hours. The time that elapsed between minima in the last observation span was a little over 3 hours. We folded the LECS lightcurve with the optical modulation period and even if ≈70% of the phase was covered, we could not detect any modulation in supersoft X-ray flux. We also note that the semi-amplitude of the optical modulation is only 0.02 mag, while the X-ray count rate has a large variation.

The time scale of the phenomenon does not give very significant upper limits on the size of the obscuring region: the upper limit on the size of an obscuring clump in the nebula, assumed to be moving at v=4000 km s⁻¹ for 5000 seconds, is 2×10¹¹ cm (a small fraction of the nova shell radius at this epoch). Assuming instead some other type of phenomenon, connected with the possible orbital period, an upper limit on the size is obtained assuming the speed of light for the obscuring source: 1.5×10¹⁴ cm.

This variability is a truly puzzling phenomenon which we do not fully understand. We note that also V1494 Aql (N Aql 1999 no.2), which was observed with Chandra and also detected as a supersoft X-ray source, showed time variability in supersoft X-rays: a flare (with increase of flux by a factor 6) that lasted for about 15 minutes and pulsations every 42 minutes (Starrfield et al. 2001). Even for this nova, the supersoft X-ray variability time scale was very short.

3 SPECTRAL ANALYSIS AND INTERPRETATION

In Paper I we made the working hypothesis that the supersoft flux observed in November 1999 was entirely due to the central hot white dwarf remnant. Neglecting the LECS flux below 0.8 keV we simultaneously fitted the LECS and MECS spectra with a MEKAL model of thermal plasma (included in the software package XSPEC, see Arnaud et al. 1986) with parameters: N(H)≈2×10²¹ cm⁻², kT≈700 eV, and unabsorbed flux ≈10⁻¹² erg cm⁻² s⁻¹ (the reduced χ² was ≈1.2). In this work we analyse instead the properties of the supersoft X ray flux. The spectral distribution observed for V382 Vel, shown in Fig.2 in the range 0.1-1.1 keV, is strikingly different from the one observed with the BeppoSAX LECS for three non-nova supersoft X-ray sources: Cal 87 (Parmar et al. 1997), Cal 88 (Parmar et al. 1987) and RX J0925.7-4758 (Hartmann et al. 1999). The LECS spectra are shown in Fig. 7 of Hartmann et al. (1999) and Fig. 3 of Parmar et al. (1997). V382 Vel is much more luminous than the other observed sources, it appears to have additional, harder spectral components and above all it is very luminous also in the very soft range, 0.1-0.2 keV. It is the only object of this kind ever studied with the LECS.

We tried to fit the whole LECS and MECS spectra in the 0.1-10 keV range adding the MEKAL model used for the hard flux to a) a blackbody, b) a blackbody with absorption edges (included in XSPEC), and c) more detailed model atmospheres (see Hartmann & Heise 1997, Hartmann et al. 1999). The latter models were used to fit the non-nova supersoft X-ray sources and very reasonable results were obtained. We expected that the blackbody would not give a perfect fit (since white dwarf atmospheres resemble blackbodies only in first approximation) but we thought that it would provide an approximate estimate of the luminosity, of the range of effective temperatures, and indicate whether absorption edges must be included. Adding absorption edges
to the black-body (the most likely seemed C vi at 0.49 eV, but we also tried N vi at 0.55 keV and 0.67 keV, O viii at 0.87 keV) did not improve the fit.

We also found that a reasonable fit to the supersoft portion of the spectrum could not be obtained with any of the atmospheric models we tested, in Local Thermodynamical Equilibrium (hereafter, LTE), and in NON-LTE (hereafter, NLTE), developed by Heise et al. (1994), Hartmann & Heise (1997) and Hartmann et al. (1999). For the NLTE case, we tested four small grids of models with log(g) between 8 and 9, with four different set of abundances. The first set had cosmic abundances (developed primarily for galactic sources, hereafter NLTE-1 models), in the second set of abundances the abundances were depleted to 0.25 the solar value (developed for LMC sources, here after called NLTE-2 models), and in the third and fourth group of models the abundances were enhanced in C, N and O and in Ne, O and Mg, respectively. These models include line opacities. The two grids with enhanced abundances are still unpublished (Hartmann 2002, private communication). With either models in LTE and NLTE we could obtain a reasonable fit with the other available models and also the NLTE-2 LMC-type models was that we could not obtain a good fit by, for instance, decreasing log(g) and “tuning” the MEKAL component accordingly. In Table 1 we give the best fit parameters the NON-LTE model atmospheres with log(g)=8.5 and a MEKAL bremsstrahlung component. The value log(g)=8.5 does not give an acceptable value of χ², but it is smaller than for other values of log(g). The only reason for testing also the NLTE-2 LMC-type models was that we could not obtain a reasonable fit with the other available models and wanted to experiment with all the available grids. We concluded that no model is acceptable: with the best NLTE-1 model, model A in the Table, we obtain χ²=4.5 for 105 degrees of freedom. The fit to the data, shown in the upper part of Fig. 3, predicts counts at energies higher than 200 eV that differ from those observed by more than 5%, which is the likely uncertainty in the knowledge of the LECS spectral response. Apart from what can be interpreted as an iron line at ≈6.5 keV, at low energies (<0.2 keV) the predicted counts exceed those observed by at least 100% (see Fig.2). However, the uncertainty in the low-energy LECS response in this range is estimated not to exceed ≈20% (Parmar et al. 1997).

In a Chandra HRC-S+LETG observation of V382 Vel done in 2000 March by one of us, S. Starrfield, the HRC-S+LETG spectrum, to be described in detail in a forthcoming paper (Starrfield et al. 2002) is very different from another nova which also appeared as a supersoft X-ray source, V1494 Aql (N Aql 1999 no.2, see Starrfield et al. 2001). The main difference is the lack of conspicuous continuum for V382 Vel in March 2000: it was instead an emission line spectrum, with many high ionization emission lines in the supersoft range. These lines presumably have an origin in the ejected nebula. The BeppoSAX instruments would have detected a blend of such lines as a featureless “pseudocontinuum”. Emission lines in the supersoft range may well have existed even at the earlier epoch of our BeppoSAX observations. These lines may have been due to shock ionization in the ejecta, rather than to photoionization by the central source.

Even if the spectral resolution of the BeppoSAX LECS and MECS is not sufficient to detect lines, we tried to determine which emission lines may explain the observed X-ray spectrum and what constraints we could derive on the level of the continuum. In model C of Table 1 we added not only the bremsstrahlung continuum at the harder energy (kT=0.84 keV) and a line at ≈6.5 keV (which is necessary but is not in the range where the bulk of the flux is emitted), but also one or more “softer” spectral features in emission. As an experiment, we obtained best fits with χ²=1.6–1.7 adding such lines to both NLTE-1 and NLTE-2 models with log(g)=8.5 (see Table 1). For the NLTE-1 model, in the best fit the added gaussian feature has to be at 449 keV. It could probably be C vi (perhaps a blend of the C viii triplet at 27–28 Å). We tried to add additional lines and improved the fit only marginally, although another emission line at 243 eV may be present (probably the Fe xv line at 50.5 Å), and a narrow iron line at ≈6.4 keV is needed to explain the excess at this energy (Model C). This fit is shown in Fig. 3 (lower panel).

With the NLTE-2 model we obtained model F with χ²=1.62 adding a line (perhaps N vii) at ≈490 keV, and again an iron line for the excess at ≈6.5 keV. The total unabsorbed flux in the lines in these two models is a negligible fraction of the total bolometric flux, less than 1%. However, the flux in the line at ≈450 or 490 keV would be about 30% of the absorbed flux in the BeppoSAX LECS range. In model C, the bolometric luminosity at a distance of 2 kpc (Della Valle et al. 1999) is 6.8 × 10^{38} erg s^{-1}. This value is

### Table 1

| Model | N(H) | Teff(×10^5 K) | kT (keV) | F₂ (keV) | F₁ keV | χ²/dof | dof |
|-------|------|--------------|----------|----------|--------|--------|-----|
| A: NLTE-1+MEKAL | 1.38 | 5.6 | 0.896 | 4.53 | 105 |
| B: NLTE-1+MEKAL+L(0.449 keV) | 2.03 | 4.0 | 0.801 | 1.70 | 102 |
| C: NLTE-1+MEKAL+L(0.243 keV+0.449 keV) | 2.07 | 4.0 | 0.836 | 6.1×10^6 | 1.67 | 96 |
| D: NLTE-2+MEKAL | 1.56 | 5.3 | 0.069 | 12.66 | 106 |
| E: NLTE-2+MEKAL+L(0.453 keV) | 1.98 | 5.2 | 0.826 | 1.57 | 117 |
| F: NLTE-2+MEKAL+L(0.491 keV+0.65.2 keV) | 1.98 | 5.2 | 0.813 | 1.2×10^6 | 1.62 | 112 |
| G: NLTE/CNO+MEKAL | 1.68 | 4.8 | 0.063 | 15.81 | 106 |
| H: NLTE/NeO+Mg+MEKAL | 1.69 | 5.0 | 0.068 | 10.73 | 105 |
higher, but not much in excess of the model predictions for a $\approx 1 \, M_\odot$ WD emitting at Eddington luminosity.

The main problem in determining the white dwarf parameters accurately is the excess at low energy ($kT < 200$ eV), which cannot be fit with one or more narrow lines. Around 150 eV several transitions exist due to Fe, Si, Mg, and Ni, that could produce an intricate pattern, must be heavily absorbed, and cannot be resolved with the spectral resolution of the BeppoSAX LECS. By the time this nova was observed with the Chandra LETG, it had become much less luminous and the spectral structure had definitely changed. We only remark that a complicated multi-temperature structure most likely existed in the ejecta, and that continuum and emission lines with different origins (white dwarf for the first and nebular for the latter) can explain the complicated spectrum we detected with the BeppoSAX LECS and MECS. We rule out that our determination of effective temperature and gravity of the white dwarf can be accurate if nebular lines overlap with the white dwarf continuum. It is only clear that at this stage the atmosphere of the central source was still the dominating component of the X-ray flux.

4 CONCLUSIONS

The observation of V383 Vel 6 months after the outburst revealed a very luminous supersoft X-ray source, comparable in luminosity only with V1974 Cyg at maximum. The details of this observation raise new, unexpected and very interesting, questions. We observed irregular variability on a time scale of minutes in the supersoft flux of V382 Vel. The lack of energy dependence of the variability measured below 0.7 keV seems to rule out the ejection of an obscuring clump, yet other considerations seems to rule out orbital variability as well.

Moreover, the BeppoSAX spectrum of V382 Vel in November 1999 appears much more complex than the expected thermal continuum of a hot white dwarf plus a residual thermal component from the nebula. We cannot justify this spectrum without invoking emission lines in the supersoft range (which were indeed observed shortly after this observation with Chandra), so we suggest that the observed “supersoft X-ray source” in V382 Vel is characterized by unresolved narrow emission lines superimposed on the atmospheric continuum. We found that even a contribution of the lines of less than 1% to the total bolometric flux can significantly change the shape of the stellar continuum, and make the task of determining white dwarf temperature and effective gravity impossible with the resolution of the BeppoSAX instruments.

However, we conclude that the bulk of the X-ray flux was still due to the atmospheric continuum and not to lines at this epoch, unlike in the later observation performed with the Chandra LETG (Starrfield et al. 2002).

Emission from classical novae in outburst can be quite complex and different from one nova to another. The X-ray spectra of N LMC 1995 (Orio & Greiner 1999) and of the recurrent nova U Sco (Kahabka et al. 1999) could be fit well with model atmospheres, although U Sco required also a nebular component at higher energy. We caution however, that the spectral structure may be as complex as the one observed for Cal 83 by Paerels et al. (2001), when observed with higher resolution.

The situation was more complex for V1974 Cyg, where a model atmosphere and a hotter, thermal component were necessary to fit the spectrum. The relative importance of the two components seemed to vary in each observation (see Balman et al. 1998) and the interplay between them was rather complicated, also due to the lower energy resolution of the PSPC compared to the LECS (a factor 2.4 less) and limited energy range of the PSPC (which could not cover the harder component well, at least at the beginning). We wonder whether the spectrum of V1974 Cyg also had superimposed nebular emission lines that made it appear hotter than it was, because a lower effective temperature ($\leq 20$ eV) of the post-nova white dwarf atmosphere may explain a puzzling fact. An ionization nebula was detected only in 1994 RAS, MNRAS 000, 1–?
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Figure 2. Spectra observed in the 0.1-8 keV range with the BeppoSAX LECS and MECS in November 1999 and (above) best fit with model A (MEKAL + NLTE model atmosphere for log(g)=8.5 and cosmic abundances) and (below) with model C, which includes also three superimposed emission lines, at 0.243 keV, 0.449 keV and 6.4 keV). The first fit is not acceptable, and this is due to in a large portion to structured residuals in the range 0.1-0.2 keV. The second fit is overall much better, but the residuals in the soft range are still large. The lower panel in each plot shows the residuals, as ratio of observed over predicted counts per energy bin.

Ho for N Cyg 1992 (Casalegno et al. 2000) while other ionization lines (indicating a higher ionization potential) were not present in the nebula. We speculate therefore that in V1974 Cyg the white dwarf (Krautter et al. 1996, Balman et al. 1998) might have been cooler than it appeared by fitting the ROSAT PSPC spectrum with just a two component model.

We note that even the spectrum of a non-nova super-soft X-ray source, the Galactic MR Vel, shows non-atmospheric emission lines, attributed to a wind from the source (Bearda et al. 2002). Instead, in classical novae in outburst emission lines in the soft X-ray range could be produced by shock ionization within the nebula. Shocked, X-ray emitting material seems to be present since the beginning of the outburst (Krautter et al. 1996, Orio et al. 2001b). The BeppoSAX LECS and the ROSAT PSPC do not resolve narrow emission lines. For novae, prominent nebular emission lines in the supersoft X-ray energy range indicate interesting possibilities, specially if they should be observed with Chandra or XMM-Newton in the future. We face new questions. Are these lines at times due to collisional excitation, do shocks occurs in the nova wind even many months after the outburst? Should we consider a line driven wind at different velocity colliding into the initial radiation driven wind? The nova theory must become more detailed and refined once the X-ray spectrum is known in detail for a statistically meaningful sample of objects. The gratings in the new X-ray observatories are opening new and exciting possibilities for nova studies.

ACKNOWLEDGMENTS

We thank L. Piro, BeppoSAX Mission Scientist, for support and W. Hartmann for providing the atmospheric models he developed. M. Orio is grateful also to another colleague, C. Markwardt, for computer support and useful advice. Finally, we thank the referee, Koji Mukai. This research has been supported by the Italian Space Agency (ASI) and by the UW College of Letters & Sciences. S. Starrfield gratefully acknowledges NSF and Chandra support to ASU.

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