HNC in Protoplanetary Disks

The Harvard community has made this article openly available. Please share how this access benefits you. Your story matters.

| Citation          | Graninger, Dawn, Karin I. Öberg, Chunhua Qi, and Joel Kastner. 2015. “HNC in Protoplanetary Disks.” The Astrophysical Journal 807 [1] (July 1): L15. doi:10.1088/2041-8205/807/1/l15. |
|-------------------|-------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|
| Published Version | doi:10.1088/2041-8205/807/1/l15                                                                                                                                                                  |
| Citable link      | http://nrs.harvard.edu/urn-3:HUL.InstRepos:24820072                                                                                                                                              |
| Terms of Use      | This article was downloaded from Harvard University’s DASH repository, and is made available under the terms and conditions applicable to Open Access Policy Articles, as set forth at http://nrs.harvard.edu/urn-3:HUL.InstRepos:dash.current.terms-of-use#OAP |
HNC IN PROTOPLANETARY DISKS

DAWN GRANINGER, KARIN I. ÖBERG, CHUNHUA QI
Harvard-Smithsonian Center for Astrophysics, 60 Garden St, Cambridge, MA, 02138

AND

JIOEL KASTNER
Center for Imaging Science, School of Physics & Astronomy, and Laboratory for Multiwavelength Astrophysics, Rochester Institute of Technology, 54 Lomb Memorial Drive, Rochester NY 14623 USA

Draft version June 15, 2015

ABSTRACT

The distributions and abundances of small organics in protoplanetary disks are potentially powerful probes of disk physics and chemistry. HNC is a common probe of dense interstellar regions and the target of this study. We use the Submillimeter Array (SMA) to observe HNC 3–2 towards the protoplanetary disks around the T Tauri star TW Hya and the Herbig Ae star HD 163296. HNC is detected toward both disks, constituting the first spatially resolved observations of HNC in disks. We also present SMA observations of HCN 3–2, and IRAM 30m observations of HCN and HNC 1–0 toward HD 163296. The disk-averaged HNC/HCN emission ratio is 0.1–0.2 toward both disks. Toward TW Hya, the HNC emission is confined to a ring. The varying HNC abundance in the TW Hya disk demonstrates that HNC chemistry is strongly linked to the disk physical structure. In particular, the inner rim of the HNC ring can be explained by efficient destruction of HNC at elevated temperatures, similar to what is observed in the ISM. To realize the full potential of HNC as a disk tracer requires, however, a combination of high SNR spatially resolved observations of HNC and HCN, and disk specific HNC chemical modeling.

Keywords: Astrochemistry – ISM: molecules – Protoplanetary Disks

1. INTRODUCTION

The disks around young stars contain the gas and dust from which planets form. From theory and observations, disks are characterized by steep temperature and radiation gradients both radially and vertically. These disk structures affect planet formation (Takeuchi et al. 1996; Kretke & Lin 2012), motivating the quest for new emission-line diagnostics that can trace physical structure gradients in disks.

This study focuses on HNC, a common high density gas tracer in the ISM. At low temperatures, HNC forms efficiently together with HCN through ion-molecule reactions. HNC and HCN are observed in a variety of different sources, but their relative abundances can vary by more than order of magnitude, which has been related to the physical environment where the emission originates. In pre-stellar and protostellar environments, the observed HNC/HCN abundance ratio is ∼1 at 10 K and decreases as temperature increases (e.g. Schilke et al. 1992; Pratap et al. 1997; Liszt & Lucas 2001; Jørgensen 2004; Padovani et al. 2011; Jin et al. 2012), suggesting that thermal reaction barriers play an important role in controlling this ratio. Work by Loison et al. (2014) has shown that the HNC/HCN ratio also depends on the elemental abundances of carbon at low temperatures, due to destructive reactions with atomic carbon. In less shielded regions, the HNC/HCN ratio also depends on the radiation present (Meijerink et al. 2004).

Nearby, protoplanetary disks, such as the one surrounding TW Hya, provide an interesting test bed for evaluating the relative impact of different environmental parameters on the HNC chemistry, due to their relatively well characterized geometries and density, temperature and radiation structures. The HNC distribution in disks is interesting in its own right. HNC and HCN are two of very few nitrogen-bearing molecules detected toward disks (Dutrey et al. 1997; Chapillon et al. 2012; Obers et al. 2013), and their distributions may therefore provide some of the best available constraints on the nitrogen reservoir. Cyanides and iso-cyanides are also important for some prebiotic chemistry scenarios (Powner et al. 2004). Observing HNC in disks has proven challenging, however. The intrinsic disk scales are small and the molecular line fluxes are weak compared to interstellar clouds and protostars. To date, HNC has only been detected in a single disk, DM Tau, (Dutrey et al. 1997) and no spatially resolved observations exist.

In this letter, we report the first interferometric detection of HNC in two disks, around the T Tauri star TW Hya (D = 54 pc Qi et al. 2004) and Herbig Ae star HD 163296 (D = 122 pc Qi et al. 2011), using the SMA. We combine these observations with new and previous SMA observations of HCN toward TW Hya and HD 163296, and new IRAM 30m observations of HCN and HNC 1–0 toward HD 163296, with the aim to constrain the HNC/HCN chemistry in disks.

2. OBSERVATIONS OF HNC AND HCN

The HNC 3–2 line at 271.981 GHz was observed toward the TW Hya disk (Δα = 11°01′51.875″, Δδ =...
The same HCN line was observed toward the disk around HD 163296 ($\Delta \alpha = 17^\circ 56^m 21.28^s$, $\Delta \delta = -21^\circ 57' 22.1''$; J2000.0) on 2013 April 5 with the SMA in the compact (COM) array configuration with projected baselines of 7.7–62.1 m. J1733-130 was used as the gain calibrator. The bandpass response was calibrated using BLLAC. Flux calibration was done using observations of Neptune. The derived fluxes of J1733-130 was 1.34 Jy.

The HCN 3–2 line at 265.886 GHz was observed toward HD 163296 on 2008 April 20 and May 12 with the SMA in the compact-north (COM-N) array configuration with projected baselines of 7.8–114.3 meter. J1733-130 was used as the gain calibrator. The bandpass response was calibrated using observations of 3C273. Flux calibration was done using observations of Titan and Callisto. The derived fluxes of J1733-130 were 2.50 Jy on 2008 April 20 and 2.13 Jy on 2008 May 12.

All data were phase- and amplitude-calibrated using the MIR package. Continuum and spectral line maps were generated and CLEANed using the MIRIAD software package. Images were constructed from the TW Hya data using a Briggs (robust) weighting factor of 0 and in HD 163296 data using 2.0 and 0.3 for HCN and HNC, respectively. The resulting synthesized beam for HCN in TW Hya is $3.2'' \times 1.9''$. For HD 163296, the synthesized beams for HCN and HNC are $3.0'' \times 1.9''$ and $2.8'' \times 2.5''$, respectively.

HCN 1–0 and HNC 1–0 lines were observed as a part of a single-dish 3 mm line survey of the disk around HD 163296 on 2012 September 13 – 16 with the IRAM 30m telescope on Pico Veleta using the EMIR 90GHz receiver and the Fourier transform spectrometer (FTS) backend in average summer weather ($pwv \sim 4$–5 mm). Two spectral settings were used with the four sidebands covering 84.0 to 115.8 GHz with a resolution of 200 kHz ($\sim 0.5$ - 0.6 km/s). Pointing was checked every 1-2 hrs and focusing every 4 hrs. The data was acquired in the wobbler switching modes, with typical throws of 60-120" and an on-source integration time of 12-15 hrs per spectral setting. Spectra were baseline subtracted and averaged using CLASS. To convert from antenna temperature, $T_a^*$, to main beam temperature, $T_{mb}$, forward and beam efficiencies of 0.95 and 0.81 were applied.

3. RESULTS

3.1. Observations

HNC 3–2 emission is clearly detected and resolved toward TW Hya. Figure 1 displays the 1.1 mm continuum and HNC 3-2 SMA observations, combining the compact and extended data. The resolved integrated continuum flux is 0.78 Jy which agrees with previous studies (Qi et al. 2013). The velocity-integrated HNC flux peaks at 0.53 Jy km s$^{-1}$ beam$^{-1}$. The integrated HNC emission peaks are offset to the east and west of the continuum phase center, suggestive of a ring-like emission structure within the disk. In the HNC image the western flux peak is almost twice as strong compared to the eastern one, corresponding to $\sim 3\sigma$. Asymmetries of this magnitude may be artifacts of CLEANing process and based on visual inspection alone are agnostic about its significance. We evaluate this aspect of the the data when modeling the complex visibilities in §3.2.

To isolate the radial flux variation in the HNC emission, Fig. 1 (lower right panel) shows the deprojected azimuthally averaged HNC line flux, using an inclination angle of 6° and position angle of 155° (Qi et al. 2004, Williams & Cieza 2011). The resulting emission profile reaches a maximum at $\sim 0.5''$, indicating the presence of a small HNC hole at the center of the TW Hya disk. This agrees with the $\sim 1''$ separation between the two flux peaks in the HNC integrated emission image. The depth and size of the hole may both be affected by the image asymmetry, however. In §3.2 we evaluate the evidence for an HNC ring and quantify its location based on an analysis in the $(u,v)$-plane.

The HNC spectrum in Fig. 1 was extracted using a $27'' \times 27''$ box. The integrated line flux is 1.2 Jy km s$^{-1}$ beam$^{-1}$ ($\sim 12\sigma$). Using the the integrated HNC flux from this study and the HCN flux from Qi et al. 2008, the disk-averaged HNC/HCN emission ratio is 0.18 ± 0.06, where the error is derived using the 1$\sigma$ rms for the two observations.

HNC is also detected toward HD 163296, but with a lower SNR. Figure 2 displays the 1.1 mm continuum, HNC, and HCN emission toward HD 163296 as well as the spectra of the molecules and the velocity structure of
HNC and HCN. The velocity integrated emission maps yields a maximum HNC and HCN flux of 0.71 Jy km s\(^{-1}\) beam\(^{-1}\) and 2.71 Jy km s\(^{-1}\) beam\(^{-1}\), respectively. The HNC emission toward the HD 163296 disk is suggestive of a ring viewed at relatively high inclination, but the low SNR inhibits any further spatial analysis.

Spectra were extracted using an elliptical mask, which increased the SNR compared to a square mask. The total integrated flux from these spectra is 0.66 Jy km/s beam\(^{-1}\) (∼4σ) and 4.47 Jy km/s beam\(^{-1}\) (∼9.5σ) for HNC and HCN, respectively. This yields an HNC/HCN ratio of 0.14 ± 0.05, similar to TW Hya.

3.2. HNC line modeling toward TW Hya

To constrain the spatial extent of the HNC emission in TW Hya, we follow the procedure outlined in Qi et al. (2008); Oberg et al. (2012, 2013) to extract molecular column density profiles. First, parametric HNC abundance models are defined with respect to the TW Hya disk temperature and density structures from Qi et al. (2013b). The line emission is modeled using the two-dimensional non-LTE Monte Carlo radiative transfer code RATRAN (Hogerheijde & van der Tak 2000) using collisional cross sections from Dumouchel et al. (2010). The HNC abundance models follow the form in Qi et al. (2008), which assumes a molecular layer with a vertically constant abundance, bound by surface (\(\sigma_s\)) and midplane (\(\sigma_m\)) boundaries defined with respect to the vertically integrated hydrogen column in the unit of \(1.59 \times 10^{21}\) cm\(^{-2}\). We set the HNC boundaries using the values from Qi et al. (2008) for HCN (log\(\sigma_s\)=-0.5, log\(\sigma_m\)=1.5). The
The radial distribution of HNC was modeled using an abundance power-law with an outer radius (R_o), hole radius (R_h), and power-law index of -1.5. The parameter R_o was varied between 35 and 60 AU and R_h was varied between 20 and 50 AU. The minimum $\chi^2$ value was calculated from the modeled and observed HNC emission in the $(u,v)$-plane using the data from the compact and extended configuration.

Figure 4 shows examples of emission models and observations in the extended configuration. The best fit model based on the $\chi^2$ value is a 10 AU wide ring with a R_o of 45 AU and R_h of 35 AU (Figure 5), shown in the lower left panel of Figure 4. This model results in no significant residuals when subtracted from the data in the $(u,v)$-plane and then imaged. The observed HNC azimuthal asymmetry toward TW Hya is thus not significant. The presence of a HNC ring in the TW Hya disk is by contrast a statistically significant result. Figure 5 shows the models that fit the data within 2 and 3$\sigma$. Within 2$\sigma$, the hole size is 28–45 AU (3$\sigma$: 26–47 AU) and the outer edge is at 37–55 AU (3$\sigma$: <57 AU).

4. DISCUSSION

In this study, we detect and spatially resolve HNC emission toward TW Hya and HD 163296. From these observations we extracted disk averaged HNC/HCN emission ratios, and the inner and outer HNC emission radii in the TW Hya disk. Here, we will discuss the constraints that these observations provide on the HNC chemistry and the disk structure.

4.1. HNC/HCN Emission Ratios

The disk-averaged HNC/HCN emission ratio is 0.1 – 0.2 toward both TW Hya and HD 163296. Based on ISM observations (Schilke et al. 1992) and astrochemistry models (Graninger et al. 2014), the HNC/HCN ratio can be used to constrain the characteristic temperature of the emission region because of efficient destruction of HNC by hydrogen atoms at temperatures >25 K. However, to apply this result straightforwardly requires...
ratios are thus not useful probes of disk properties. the complete HNC/HCN chemistry, disk-averaged line
sence of detailed disk chemistry models that incorporate
C/O ratio, which may be elevated in disks. In the ab-
also shown that the HNC/HCN ratio depends on the
disks (Meijerink et al. 2007). Loison et al. (2014) has
cloud cores, but may not be a good approximation in
glects the influence of radiation, which is reasonable in
cloud cores, but may not be a good approximation in
disks (Meijerink et al. 2007). Loison et al. (2014) has
also shown that the HNC/HCN ratio depends on the
C/O ratio, which may be elevated in disks. In the ab-
presence of detailed disk chemistry models that incorporate
the complete HNC/HCN chemistry, disk-averaged line
ratios are thus not useful probes of disk properties.

Figure 5. $\chi^2$ contour plot for the best fit models. The contours are 2 and 3σ. The 3σ contour is in red and the red cross marks
the best-fit model. The black crosses mark the values of $R_h$ and $R_o$ modeled and the shaded region indicates where $R_h \geq R_o$.

that HNC and HCN are co-spatial, that there is a sim-
ple relationship between emission and abundance ratios
of HNC and HCN, and that temperature dominates the
chemistry, none of which clearly hold in the case of these
disks.

First, HNC and HCN are not co-spatial – some
HCN is present toward the center of the disk, where
HNC is absent. This may be partially mitigated by
averaging the HNC and HCN emission only across
disk areas where HCN and HNC emission are respec-
tively detected. Second, the relationship between line
emission and abundance ratios is complicated by high
HCN line opacities and different HCN and HNC exci-
tation properties Sarrasin et al. (2010). We used the
large velocity gradient radiative transfer code, RADEX
van der Tak et al. (2007), with the collisional rate con-
stants from Dumouchel et al. (2011) to evaluate relation-
ship between the HNC/HCN line emission and abun-
dance ratios using temperatures and densities character-
istic of HCN disk emission Walsh et al. (2010) and the
HCN column density from Qi et al. (2008). We find that
the line and abundance ratios can differ by up to an or-
magnitude, demonstrating that more detailed modeling
is required to extract HNC/HCN abundance ratios from
the observed line emission.

Even if abundance ratios are successfully extracted, it
may be difficult to infer a characteristic emission region
temperature. Other environmental characteristics than
temperature may regulate the HNC/HCN abundance ra-
tio in disks. The model by Graninger et al. (2014) ne-
glects the influence of radiation, which is reasonable in
cloud cores, but may not be a good approximation in
disks Meijerink et al. 2007. Loison et al. (2014) has
also shown that the HNC/HCN ratio depends on the
C/O ratio, which may be elevated in disks. In the ab-
sence of detailed disk chemistry models that incorporate
the complete HNC/HCN chemistry, disk-averaged line
ratios are thus not useful probes of disk properties.

4.2. The HNC inner and outer radii in the TW Hya
disk

In the disk toward TW Hya, there is a lack of HNC
emission interior to 28–45 AU, corresponding to tempera-
tures above 21–13 K using the Qi et al. (2013b) tem-
perature profile. HCN does not have an emission hole
Qi et al. (2008). The observed absence of HNC close to
the star is consistent with temperature dependent HNC
destruction Graninger et al. (2014). It is difficult to ex-
plain the hole with the increased UV flux present at
smaller radii, since HNC and HCN are thought to have
similar photodissociation cross-sections Roberge et al.
1991, van Dishoeck et al. 2006.

HNC also has an outer emission radius between 35–55
AU, with the best fit at 45 AU, substantially smaller than
the observed radii of other gases, including CO, HCO+
and HCN Qi et al. (2008). This outer cutoff in HNC
emission was unexpected. It coincides with the onset
of C$_2$H emission, which is proposed to be a UV pho-
tochemistry product Kastner et al. (2013), and is close
to the radius where the dust density sharply decreases
Andrews et al. (2012). These spatial coincidences sug-
gest three potential explanations for the outer cutoff
of HNC emission: 1. UV photodissociation of HNC due
to deeper penetration of stellar and interstellar UV photons
in the low-column density outer disk, 2. an increased UV
photodissociation of H$_2$ and carbon-bearing molecules
accelerating the HNC+H and HNC+C destruction pro-
cesses at a given temperature Graninger et al. 2014,
Loison et al. (2014), and 3. grain-gas thermal decoupling
due to lower gas and grain densities, causing a temper-
ature reversal in the outer disk and thus the increase in
thermally controlled HNC destruction (i.e. a similar
process that sets the inner edge).

The first two explanations both depend on a more
efficient UV photochemistry exterior to 45 AU, consist-
tent with the onset of C$_2$H emission at the same ra-
dius Kastner et al. (2013), and signs of efficient UV pho-
tochemistry in the outer low-density regions of other
disks Oberg et al. (2011). Flared disks can intercept
stellar UV radiation out to larger radii, and Ly emission
is efficiently down-scattered through disk towards
the midplane Fogel et al. (2011). Additionally, the in-
terstellar radiation field will contribute a non-negligible
amount of UV radiation to the outer disk Bergin et al.
2003. Efficient UV penetration at large disk radii may
increase the HNC photodissociation rate, reducing its
abundance. This scenario could be tested by compar-
ing HNC and HCN emission profiles. HNC and HCN are
believed to have similar photodissociation cross-sections
Roberge et al. 1991, van Dishoeck et al. 2006, and if
their outer emission radii are UV controlled, they should
be the same. The best-fit outer HCN emission radius is
~100 AU Qi et al. (2008), twice that of HNC, but this
value is not well enough constrained to exclude direct
photodissociation of HCN as an explanation for the HNC
outer radius.

The second and third proposed explanations for the
HNC outer cut-off should produce a smaller HNC than
HCN emission radius, due to an increased HCN-to-HCN
conversion rate in the outer disk. In the second scenario
and increased UV penetration increases the atomic H
and C abundance through H$_2$ and CO photodissociation

Graninger et al. (Akimkin et al. 2013), speeding up HNC isomerization at a given temperature (Graninger et al. 2014). In the third scenario, a lower density in the outer disk reduces gas-grain thermal coupling sufficiently that direct gas heating by the interstellar radiation field becomes important. This heating could result in a radial temperature gradient reversal outside of 45 AU, again increasing the HNC isomerization rate.

Higher spatial resolution observations with ALMA should be able to distinguish between the three scenarios. The relative sizes of the HNC and HCN outer radii should differentiate between #1 on one hand and #2 and #3 on the other. If #1 is ruled out, the radial temperature profile of HCN beyond 45 AU should be able to differentiate between #2 and #3. Whichever scenario is confirmed, spatially resolved HNC/HCN emission modeled using non-LTE techniques clearly has a large potential to probe interesting disk properties.

This work has benefited from discussions with Sean Andrews, David Wilner, Ryan Loomis, Eric Herbst and the helpful comments from the anonymous referee. The Submillimeter Array is a joint project between the Smithsonian Astrophysical Observatory and the Academia Sinica Institute of Astronomy and Astrophysics and is funded by the Smithsonian Institution and the Academia Sinica. The authors wish to recognize and acknowledge the very significant cultural role and reverence that the summit of Mauna Kea has always had within the indigenous Hawaiian community. We are most fortunate to have the opportunity to conduct observations from this mountain. The study is also based on observations with the IRAM 30m Telescope. IRAM is supported by INSU/CNRS (France), MPG (Germany) and IGN (Spain). KIÖ acknowledges funding from the Simons Collaboration on the Origins of Life Investigator award #321183, the Alfred P. Sloan Foundation, and the David and Lucile Packard Foundation. JK’s research on nearby, irradiated disks is supported by National Science Foundation grant AST-1108950 to RIT.

REFERENCES

Aikawa, Y., Miyama, S. M., Nakano, T., & Umemayashi, T. 1996, ApJ, 467, 684
Akimkin, V., Zhukovska, S., Wiebe, D., et al. 2013, ApJ, 766, 8
Andrews, S. M., Wilner, D. J., Hughes, A. M., et al. 2012, ApJ, 744, 162
Bergin, E., Calvet, N., D’Alessio, P., & Herczeg, G. J. 2003, ApJL, 591, L159
Chapillon, E., Dutrey, A., Guilloteau, S., et al. 2012, ApJ, 756, 58
Dumouchel, F., Faure, A., & Lique, F. 2010, MNRAS, 406, 2488
Dutrey, A., Guilloteau, S., & Guelin, M. 1997, A&A, 317, L55
Dutrey, A., Henning, T., Guilloteau, S., et al. 2007, A&A, 464, 615
Fogel, J. K. J., Bethell, T. J., Bergin, E. A., Calvet, N., & Semenov, D. 2011, ApJ, 726, 29
Graninger, D. M., Herbst, E., Öberg, K. I., & Vasyunin, A. I. 2014, ApJ, 787, 74
Hogerheijde, M. R. & van der Tak, F. F. S. 2000, A&A, 362, 697
Jin, M., Lee, J.-E., & Kim, K.-T. 2015, ArXiv e-prints
Jørgensen, J. K. 2004, A&A, 424, 589
Kastner, J. H., Qi, C., Gorti, U., et al. 2015, ArXiv e-prints
Kretke, K. A. & Lin, D. N. C. 2012, ApJ, 755, 74
Liszt, H. & Lucas, R. 2001, A&A, 370, 576
Loison, J.-C., Wakelam, V., & Hickson, K. M. 2014, MNRAS, 443, 398
Meijerink, R., Spaans, M., & Israel, F. P. 2007, A&A, 461, 793
Öberg, K. I., Guzmán, V. V., Furuya, K., et al. 2015, Nature, 520, 198
Öberg, K. I., Qi, C., Fogel, J. K. J., et al. 2011, ApJ, 734, 98
Öberg, K. I., Qi, C., Wilner, D. J., & Hogerheijde, M. R. 2012, ApJ, 749, 162
Padovani, M., Walmsley, C. M., Tafalla, M., Hily-Blant, P., & Pineau Des Forêts, G. 2011, A&A, 534, A77
Povich, M. W., Gerland, B., & Sutherland, J. D. 2009, Nature, 459, 239
Pratap, P., Dickens, J. E., Snell, R. L., et al. 1997, ApJ, 486, 862
Qi, C., D’Alessio, P., Öberg, K. I., et al. 2011, ApJ, 740, 84
Qi, C., Ho, P. T. P., Wilner, D. J., et al. 2004, ApJL, 616, L11
Qi, C., Öberg, K. I., & Wilner, D. J. 2013a, ApJ, 765, 34
Qi, C., Öberg, K. I., Wilner, D. J., et al. 2013b, Science, 341, 630
Qi, C., Wilner, D. J., Aikawa, Y., Blake, G. A., & Hogerheijde, M. R. 2008, ApJ, 681, 1396
Roberge, W. G., Jones, D., Lepp, S., & Dalgarno, A. 1991, ApJS, 77, 287
Rosenfeld, K. A., Qi, C., Andrews, S. M., et al. 2012, ApJ, 757, 129
Sarussi, A., Abdallah, D. B., Wernli, M., et al. 2010, MNRAS, 404, 518
Schilke, P., Walmsley, C. M., Pineau Des Forets, G., et al. 1992, A&A, 256, 595
Takeuchi, T., Miyama, S. M., & Lin, D. N. C. 1996, ApJ, 460, 832
van der Tak, F. F. S., Black, J. H., Schöier, F. L., Jansen, D. J., & van Dishoeck, E. F. 2007, A&A, 468, 627
van Dishoeck, E. F., Jonkheid, B., & van Hemert, M. C. 2006, Faraday Discussions, 133, 231
Walsh, C., Millar, T. J., & Niswara, H. 2010, ApJ, 722, 1607
Williams, J. P. & Cieza, L. A. 2011, ARA&A, 49, 67