ABSOLUTE PROPERTIES OF THE(binary) SYSTEM BB PEGASI

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ABSTRACT

We present ground-based photometry of the low-temperature contact binary BB Peg. We collected all the times of mid-eclipse available in the literature and combined them with those obtained in this study. Analyses of the data indicate a period increase of $(3.0 \pm 0.1) \times 10^{-8}$ days yr$^{-1}$. This period increase of BB Peg can be interpreted in terms of the mass transfer $2.4 \times 10^{-8}$ M$_\odot$ yr$^{-1}$ from the less massive to the more massive component. The physical parameters have been determined as $M_c = 1.42$ M$_\odot$, $M_b = 0.53$ M$_\odot$, $R_c = 1.29$ R$_\odot$, $R_b = 0.83$ R$_\odot$, $L_c = 1.86 L_\odot$, and $L_b = 0.94 L_\odot$ through simultaneous solutions of light and of the radial velocity curves. The orbital parameters of the third body, which orbits the contact system in an eccentric orbit, were obtained from the period variation analysis. The system is compared to the similar binaries in the Hertzsprung-Russell and mass-radius diagrams.

Key words: binaries: close — binaries: eclipsing — stars: individual (BB Pegasi) — stars: late-type

Online material: color figures, machine-readable table

1. INTRODUCTION

BB Peg (HIP 110493; $V = 11.6$ mag, F8 V) is a low-temperature contact binary (LTCB) system which was discovered as a variable star in 1931 by Hoffmeister (1931). Whitney (1959) refined the orbital period. Since then BB Peg has been the subject of several investigations. The system was observed photoelectrically in 1978 by Cerruti-Sola & Scaltriti (1980), Zhai & Zhang (1979), and Awadalla (1988). The times of minima of the system have been published by numerous authors.

Cerruti-Sola et al. (1981) analyzed the $BV$ light curves of Cerruti-Sola & Scaltriti (1980) using the Wilson-Devinney (WD; Wilson & Devinney 1971) code. Giuricin et al. (1981) solved the same light curves using the Wood (1972) model and obtained somewhat different results. Leung et al. (1985) used WD to analyze the $BV$ light curves obtained by Zhai & Zhang (1979). Awadalla (1988) observed $UBV$ light curves but did not perform a light-curve analysis. The mass ratio was determined photometrically for these light-curve solutions. The first radial velocity study of the system done by Hrivnak (1990) gives the mass ratio as 0.34 ± 0.02. More recent radial velocity data obtained by Lu & Rucinski (1999) result in a mass ratio of 0.360 ± 0.006. The photometric mass ratio $(0.360 \pm 0.003)$ derived by Leung et al. (1985) agrees very well with the spectroscopic value, a result of the total/annular nature of the eclipses (see Terrell & Wilson 2005). Zola et al. (2005) published the physical parameters of the components. The orbital period variation was studied by Cerruti-Sola & Scaltriti (1980) and Qian (2001).

2. OBSERVATIONS

The photometric observations of the system were obtained with the 0.4 m (T40), 0.35 m (T35), and 0.30 m telescopes (T30) at the Ege University Observatory and TÜBİTAK National Observatory on eight nights during the observing season between August and December 2004 with T35 and two nights in 2006 with T40. However, the system was observed at T30 and T40 for only three nights in order to obtain the minimum times. The light curve of the system was obtained from CCD photometry observations. The light curves of BB Peg in the Bessel $V$ and $R$ filters are shown in Figure 2d (discussed in § 4), and the data are given in Table 1. The comparison and check stars were BD +15 4634 and GSC 01682–01530, respectively.

We obtained two minimum times throughout these observations. They are listed in Table 2, together with those published in the existing literature. Using these minimum times we derived the linear ephemeris

$$HJD \text{ Min. } I = 2.450.657.4599(4) + 0.3615015(1)E \quad (1)$$

and used it in the reduction processes of the observed data.

3. ECLIPSE TIMINGS AND PERIOD STUDY

The period variation study of the system was presented for the first time by Cerruti-Sola & Scaltriti (1980), resulting in the ephemeris Min. I ($HJD) = 2.443.764.3334(6) + 0.3615021(2)E + (2.3 \times 10^{-11})E^2$. Qian (2001) presented it as Min. I ($HJD) = 2.430.285.7618(6) + 0.36150027(1)E + [(2.35 \pm 0.01) \times 10^{-11}]E^2$.

Recently, the existence of a third body was reported via spectroscopic study by D’Angelo et al. (2006). We used the linear ephemeris given by Qian (2001) to construct the binary’s $O - C$ diagram. It shows almost a sine-like variation superposed on an upward parabola. A sine-like variation in the $O - C$ curve, where both the primary and the secondary minima follow the same trend, suggests the light-time effect via the presence of a tertiary component. Times of minima of BB Peg yielded the following equation:

$$\text{Min. } I = T_0 + P_0E + \frac{1}{2} \frac{dP}{dE} E^2 + \frac{a_{12} \sin i'}{c} \times \left[ \frac{1 - e'^2}{1 + e' \cos i'} \sin (i' + \omega') + e' \sin \omega' \right]. \quad (2)$$

where $T_0$ is the starting epoch for the primary minimum; $E$ is the integer eclipse cycle number; $P_0$ is the orbital period of the
### Table 1

**VR Measurements of BB Peg**

| HJD   | Phase | Δm   | Filter |
|-------|-------|------|--------|
| 53,301.3054 | 0.5115 | 1.2930 | 1     |
| 53,301.3067 | 0.5151 | 1.2840 | 1     |
| 53,301.3080 | 0.5187 | 1.2610 | 1     |
| 53,301.3093 | 0.5224 | 1.2500 | 1     |
| 53,301.3106 | 0.5260 | 1.2300 | 1     |
| 53,301.3119 | 0.5296 | 1.2050 | 1     |
| 53,301.3133 | 0.5333 | 1.1840 | 1     |
| 53,301.3146 | 0.5369 | 1.1690 | 1     |
| 53,301.3159 | 0.5405 | 1.1560 | 1     |

**Notes.**—See Fig. 2d. The phases were calculated using eq. (1). In the fourth column, 1 and 2 denote the V and R filters, respectively. Table 1 is published in its entirety in the electronic edition of the Astronomical Journal. A portion is shown here for guidance regarding its form and content.

### Table 2

**The Primary (I) and Secondary (II) Minima Times in HJD**

| HJD       | Min. | Ref. | HJD       | Min. | Ref. | HJD       | Min. | Ref. |
|-----------|------|------|-----------|------|------|-----------|------|------|
| 26,559.241 | II   | 1    | 41,181,397 | I    | 5    | 50,657,4575 | I   | 16   |
| 26,582.014 | II   | 2    | 41,335,227 | I    | 6    | 50,671,3770 | I   | 16   |
| 26,965.204 | II   | 2    | 42,405,259 | I    | 7    | 50,702,4698 | I   | 17   |
| 27,393.223 | II   | 2    | 42,607,523 | I    | 8    | 50,739,7052 | II  | 18   |
| 30,226.826 | I    | 3    | 42,748,310 | II   | 8    | 50,769,525  | I   | 18   |
| 30,235.865 | I    | 3    | 42,729,449 | I    | 9    | 51,078,4304 | II  | 19   |
| 30,258.638 | I    | 3    | 43,730,351 | I    | 9    | 51,471,3810 | I   | 20   |
| 30,264.797 | I    | 3    | 43,754,389 | I    | 9    | 52,131,8425 | II  | 21   |
| 30,281.776 | I    | 3    | 43,754,389 | I    | 9    | 52,201,2508 | I   | 22   |
| 30,285.753 | I    | 3    | 43,757,466 | I    | 9    | 52,201,4305 | I   | 22   |
| 30,530.861 | I    | 3    | 43,764,333 | I    | 9    | 52,203,2386 | I   | 22   |
| 30,552.903 | I    | 3    | 43,806,0845 | I  | 10   | 52,203,4188 | I   | 22   |
| 30,584.721 | I    | 3    | 43,806,9883 | I  | 10   | 52,207,3962 | I   | 22   |
| 30,994.128 | I    | 2    | 43,813,1336 | I  | 10   | 52,513,4118 | I   | 23   |
| 31,731.756 | I    | 4    | 43,814,0371 | I  | 10   | 52,838,402  | I   | 24   |
| 31,783.455 | I    | 4    | 43,842,0537 | I  | 10   | 52,852,4956 | I   | 25   |
| 32,433.631 | II   | 4    | 43,866,9989 | I  | 10   | 52,243,4607 | II  | 26   |
| 32,433.801 | I    | 4    | 44,812,503 | II   | 11   | 52,284,3112 | II  | 22   |
| 32,436.687 | I    | 4    | 45,208,351 | I    | 12   | 52,285,3957 | II  | 22   |
| 32,436.866 | II   | 4    | 45,208,531 | II   | 12   | 52,353,3577 | II  | 26   |
| 32,453.697 | II   | 4    | 46,024,2600 | II  | 12   | 53,984,3589 | I   | 26   |
| 32,455.683 | II   | 4    | 46,026,2483 | I    | 12   | 53,984,3591 | I   | 26   |
| 32,473.567 | I    | 4    | 49,243,4462 | II   | 13   | 53,984,5409 | II  | 26   |
| 32,477.538 | II   | 4    | 49,244,3490 | I    | 13   | 53,984,5411 | II  | 26   |
| 32,477.744 | II   | 4    | 49,273,2689 | I    | 13   | 53,986,5271 | I   | 26   |
| 32,479.710 | II   | 4    | 49,275,2600 | II   | 13   | 53,986,5276 | II  | 26   |
| 34,711.615 | I    | 4    | 50,001,3351 | I    | 14   | 53,992,4949 | I   | 26   |
| 35,468.604 | I    | 4    | 50,026,2785 | I    | 14   | 53,992,4957 | I   | 26   |
| 36,056.764 | I    | 4    | 50,359,4028 | II   | 15   | 26   |

**Note.**—HJD = 2,400,000.

**References.**—(1) Zesseteitach 1939; (2) Tsessevisch 1954; (3) Whitney 1943; (4) Whitney 1943; (5) Diehelm 1973; (6) Locher 1973; (7) Diehelm 1976; (8) Diehelm 1977; (9) Cerruti-Sola & Scaltriti 1980; (10) Zhai & Zhang 1979; (11) Derman et al. 1982; (12) Awadalla 1988; (13) Müyesseroglu et al. 1996; (14) Agerer & Hübscher 1996; (15) Agerer & Hübscher 1998a; (16) Ogloza 1997; (17) Agerer & Hübscher 1998b; (18) Samolyk 1999; (19) Agerer et al. 1999; (20) Agerer et al. 2001; (21) Nelson 2002; (22) Drozd & Ogloza 2005; (23) Demircan et al. 2003; (24) Batık et al. 2003; (25) Hübscher 2005; (26) this work.

### 4. Light-Curve Analysis

Previous light curves have been analyzed either by old methods or with the assumption that the photometric mass ratio was known. All previously published light curves, as well as those of the current study, have been analyzed simultaneously with the Lu & Rucinski (1999) radial velocities using the latest version of the WD code (Wilson & Devinney 1971; Wilson 1994). Mode 3 of the WD code has been used throughout the analysis. As seen in Figure 2, the light curves show asymmetries in the maxima. Generally, it is accepted that stellar activity may cause these asymmetries in the light curves; we discuss these asymmetries in § 5. Hence, the stellar spot parameters were taken into consideration.
in our analysis. The adopted values are $T_1 = 6250$ K and, according to the $B - V$ color index, gravity-darkening coefficients and albedos were chosen as $g_1 = g_2 = 0.32$ (Lucy 1967) and $A_1 = A_2 = 0.5$ (Rucinski 1969) and the logarithmic limb-darkening coefficients ($x_1$, $x_2$) were obtained from van Hamme (1993). The semimajor axis of the relative orbit $a$, binary center-of-mass radial velocity $V_{c13}$, inclination $i$, temperature of the secondary component $T_2$, luminosities of the primary component $L_1$ ($U$, $B$, $V$, $R$), potential of the common surface $\Omega$, and spot parameters (latitude, longitude, size, and temperature factor) were adjustable parameters. The results are given in Table 4. Weights for the different sets of data were determined by the scatter of the observations. In all the analyses the $B$, $V$, and $R$ filters were given 2 times higher weight than the $U$ filter to take their much better dispersion into account. The computed light curves (solid lines) obtained along with the parameters given in Table 4 were compared with all the observed light curves shown in Figures 2a–2d. The synthetic light curves were created with the LC program.

The obtained parameters for the light curves are given in Table 4. The results of the different light-curve solution models (M) have been denoted by different numbers. We have assigned M1 in Table 4 to two colors ($B$ and $V$) with light-curve solutions obtained from Cerruti-Sola & Scaltriti (1980), M2 to two colors ($B$ and $V$) with the light-curve model of Zhai & Zhang (1979; the mean values are taken from Leung et al. 1985), M3 to three colors ($U$, $B$, and $V$) with light curves that were obtained by

### Table 3

**Orbital Elements of the Tertiary Component in BB Peg**

| Parameter                  | Value                      |
|----------------------------|----------------------------|
| $T_0$ (HJD)                | 2,430.285.7655(36)         |
| $P_0$ (days)               | 0.3615006(1)               |
| $P_0$ (yr)                 | 27.9(2.0)                  |
| $T^*$ (HJD)                | 2,438.540(793)             |
| $e'$                       | 0.56(0.30)                 |
| $\omega'$ (deg)            | 69(18)                     |
| $a_{13} \sin i'$ (AU)      | 0.96(15)                   |
| $f(m) (M_\odot)$           | 0.0010(5)                  |
| $m_{3.5, i = 0^\circ} (M_\odot)$ | 1.23                    |
| $m_{3.5, i = 90^\circ} (M_\odot)$ | 0.16                    |
| $Q$ (counts day$^{-1}$)    | $1.5(2) \times 10^{-11}$  |

**Note:** The standard errors (1σ) in the last digit are given in parentheses.
Awadalla (1988), and M4 to two colors (R and V) with light curves obtained in this study. All the results appear to be compatible with each other. Consistency of observations, using the results given in Table 4, with applied models is shown in Figures 2a–2d.

Keeping in mind the possibility of a tertiary component orbiting a third body orbiting the binary system, we assume the third body’s (l3) parameter as a free parameter through the light-curve solution. However, we could not find meaningful values for the l3 parameter throughout the solutions. Likewise, D’Angelo et al. (2006) showed that the light contribution of the third body is tiny (l3/l1+2 = 0.009).

5. RESULTS AND DISCUSSION

All available light curves in the literature have been solved using the recent WD code, and the results are presented in Table 4. The solutions yielded very similar results. During the process the effective temperature and absolute magnitude of the Sun were taken as 5780 K and 4.75 mag, respectively. In Figure 3 the component parameters are shown on the H-R and mass-radius dia-

grams. We show them along with the LTCB systems (Yakut & Eggleton 2005), whose physical parameters are well known. The results obtained from analyzing BB Peg (Table 5) seem to be in good agreement with the well-known LTCBs. The location of the less massive component in the system indicates that the system is overluminous and oversized, like the other W-subtype secondary stars. Companion stars appear to be below the zero-age main sequence, and the massive component is situated near the terminal-age main sequence. If interstellar absorption is not taken into account, then through the parameters given and using the values given in Table 5 the distance of the system is found to be 361 ± 25 pc. This is consistent with the Hipparcos value (Perryman et al. 1997). The system’s distance is derived from the Rucinski & Duerbeck (1997) period-color-luminosity relation 389 pc, which is close to that obtained in this study.

Many contact binaries show an asymmetry in which one maximum is higher than the other (the O’Connell effect); these asymmetries are usually attributed to spots, which we interpret here in a very general sense: they might be due to large cool starspots, to hot regions such as faculae, to gas streams and their
impact on the companion star, or to some inhomogeneities not yet understood (Yakut & Eggleton 2005). While the asymmetry is apparent in the shape of the light curve of some systems (e.g., YZ Phe; Samec & Terrell 1995), in others this asymmetry may not be so prominent (e.g., XY Leo; Yakut et al. 2003). The asymmetry in the light curve of BB Peg is modeled with a cold spot on the secondary component (the cooler component with higher mass and radius) of the system. In the model of the light curve denoted by M2 the spot activity appears to be prominent with respect to the other models. The results of the model are summarized in Table 4. In addition, the asymmetry in the light curve is well represented by the model (see Fig. 2).

Figure 1 shows a parabolic variation. Therefore, we have applied a parabolic fit and assume that the mass transfer takes place between the components. The parabolic \((\Delta T_1)\) curve shown in Figure 1 indicates the existence of mass transfer in the contact system BB Peg. Upward parabolic variation suggests mass transfer from the less massive component (the hotter component in the case of BB Peg) to the more massive component. Equation (2) yields a period increase at a rate of \(dP/dt = (3.0 \pm 0.1) \times 10^{-8}\) days yr\(^{-1}\). If the period increase is indeed caused by conservative mass transfer, then one can estimate the mass transfer between the components. Using the derived masses, we derive the rate of mass transfer, \((2.4 \pm 0.4) \times 10^{-8}\) \(M_\odot\) yr\(^{-1}\), from the less massive to the more massive component as in the conservative mass transfer approximation. However, conservative mass transfer is just an optimistic assumption. The non-conservative case is very important in close binary evolution (for details, see Yakut & Eggleton 2005 and references therein). Analysis of the data, obtained over approximately 25 years, using the WD program indicates a period increase of \((2.9 \pm 0.1) \times 10^{-8}\) days yr\(^{-1}\), which is close to that obtained with \(O - C\) analysis.

Figure 1b shows the variation of \(\Delta T_II\) when the observations are extracted from the parabolic variation. The \(\Delta T_II\) variations show a sine-like variation, which implies the existence of a tertiary component orbiting BB Peg on an eccentric orbit. Spectroscopic study of the system shows the existence of an M-type dwarf star about the binary (D’Angelo et al. 2006). Using this information with sine-line variation of the residuals of \((O - C)\), we solved the system under the assumption of the existence of a third body and obtained the values given in Table 3. The results of the \((O - C)\) analysis show that the third component has a highly eccentric orbit \((e = 0.56)\) with about a 30 yr period. Indeed, \((O - C)\) residuals may indicate that the source of this variation could be due to magnetic activity. The orbit of the third body obtained in this study, compared to the data of Pribulla & Rucinski (2006), appeared to be much more eccentric.

![Fig. 3.—H-R and mass-radius diagrams showing BB Peg. The filled circles show the primary component of W-type LTCBs, and the open circles represent the secondaries. The zero-age main sequence line is taken from Pols et al. (1995). [See the electronic edition of the Journal for a color version of this figure.]](image)

| TABLE 5 | Absolute Parameters of BB Peg |
|----------------|-----------------------------|
| Parameter                  | Hot Component | Cool Component |
| Mass \((M_\odot)\)      | 0.53(2)        | 1.42(4)         |
| Radius \((R_\odot)\)    | 0.83(2)        | 1.29(2)         |
| Effective temperature \((K)\) | 6250          | 5950(30)        |
| Luminosity \((L_\odot)\) | 0.94(6)        | 1.86(8)         |
| Surface gravity (cgs)     | 4.33           | 4.37            |
| Absolute bolometric magnitude (mag) | 4.82\(^{+0.08}_{-0.09}\) | 4.09\(^{+0.13}_{-0.08}\) |
| Absolute visual magnitude (mag) | 4.98          | 4.26            |
| Distance (pc)             | 361\(^{+25}_{-20}\) | 361\(^{+25}_{-20}\) |

Note.—The standard errors \((1 \sigma)\) in the last digit are given in parentheses.
On the other hand, using the values given in Tables 3 and 5 one may predict the mass of the tertiary component. By assigning 0.96 AU to $a_{12}\sin i$ and 29.7 yr as the period, one can give the mass function as $0.0010 M_{\odot}$. For orbital inclinations ($i_2$) of 90, 80, 50, 30, and 10 the masses of the third body ($m_3$) are estimated to be 0.161, 0.164, 0.214, 0.341, and 1.229 $M_{\odot}$, respectively. D’Angelo et al. (2006) found a temperature of 3900 K for the tertiary component and a luminosity ratio $\left[\frac{\beta}{\lambda_3}\right] = \frac{l_3}{l_1 + l_2}$ of 0.009. Following this information with the deduced luminosities given in this study, one may give the radius of the third body as $0.33 R_{\odot}$. The $M'_{0.978}$ relationship is deduced using the 10 well-known M-type dwarf stars given in the study of López-Morales & Ribas (2005), then the tertiary body’s mass of 0.32 $M_{\odot}$ is found. Taking into consideration that value of mass, the orbital inclination of the third body is found to be 35°. Useful observations of BB Peg throughout the next decade will help to determine the accurate orbital parameters of the third body from the $O - C$ diagram.

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