ALMA Detection of Extended [C II] Emission in Himiko at $z = 6.6$

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Abstract

Himiko is one of the most luminous Ly$\alpha$ emitters at $z = 6.595$. It has three star-forming clumps detected in the rest-frame UV, with a total SFR = 20 $M_\odot$ yr$^{-1}$. We report the Atacama Large Millimeter/submillimeter Array (ALMA) detection of the [C II]158 $\mu$m line emission in this Galaxy with a significance of 8$\sigma$. The total [C II] luminosity ($L_{\text{C II}} = 1.2 \times 10^8 L_\odot$) is fully consistent with the local $L_{\text{C II}}$–SFR relation. The ALMA high-angular resolution reveals that the [C II] emission is made of two distinct components. The brightest [C II] clump is extended over 4 kpc and is located on the peak of the Ly$\alpha$ nebula, which is spatially offset by 1 kpc relative to the brightest UV clump. The second [C II] component is spatially unresolved (size < 2 kpc) and coincident with one of the three UV clumps. While the latter component is consistent with the local $L_{\text{C II}}$–SFR relation, the other components are scattered above and below the local relation. We shortly discuss the possible origin of the [C II] components and their relation with the star-forming clumps traced by the UV emission.

Key words: galaxies: evolution – galaxies: formation – galaxies: high-redshift – infrared: ISM

1. Introduction

One of the key science goals of the Atacama Large Millimeter/submillimeter Array (ALMA) is the detection and investigation of star-forming galaxies in the early universe through observations of the far-infrared (FIR) lines. In particular, ALMA observations of the strongest FIR lines, such as [C II]158 $\mu$m and [O I]88 $\mu$m, allow us to probe the properties of the interstellar medium (ISM) in galaxies at $z > 6$ (e.g., Maiolino et al. 2015; Inoue et al. 2016; Pentericci et al. 2016; Carniani et al. 2017b; Matthee et al. 2017; Smit et al. 2017), which are responsible for cosmic reionization.

The [C II] ($^2P_{3/2} \rightarrow ^2P_{3/2}$) transition at 157.74 $\mu$m, whose emissivity is mainly driven by the collisional excitation mechanism, is one of the dominant coolants of the neutral and partly ionized ISM. Furthermore, [C II] is directly correlated to star formation, since its emission can be excited by the UV radiation field of young stars in photodissociation regions. This transition can therefore be used as a tracer of ongoing star formation in local and distant star-forming galaxies (e.g., De Looze et al. 2014; Herrera-Camus et al. 2015; Vallini et al. 2015), except for ultra-luminous systems (Díaz-Santos et al. 2013; Kapala et al. 2015; Smith et al. 2017).

In the last five years ALMA observations targeting the [C II] line emission in “normal” star-forming galaxies at $z > 6$ (with SFR = 3–100 $M_\odot$ yr$^{-1}$) have yielded a variety of results (Ouchi et al. 2013; Ota et al. 2014; Willott et al. 2015; Bradač et al. 2016; Knudsen et al. 2016; Maiolino et al. 2015; Pentericci et al. 2016; Matthee et al. 2017; Smit et al. 2017).

The $L_{\text{C II}}$–SFR relation observed in these high-$z$ systems seems to have an intrinsic dispersion larger than that observed in the local universe, which probably results from the wide range of global properties characterizing primeval galaxies (Vallini et al. 2015; Carniani et al. 2017a; Lagache et al. 2017; Matthee et al. 2017). The properties of the [C II] emission are even more puzzling in luminous Ly$\alpha$ emitters (LAEs) whose detections or upper limits (Walter et al. 2012; Ouchi et al. 2013; Ota et al. 2014; Schaerer et al. 2015; Willott et al. 2015) place these sources below the $L_{\text{C II}}$–SFR relation found in the local universe (De Looze et al. 2014). So far, the only LAE having a [C II] emission consistent with the local relation is CR7 at $z \sim 7$ (Matthee et al. 2017). Simulations and models have attempted to explain the discrepancy with the local relation by taking into account the low metal content in high-$z$ galaxies, the excitation by the cosmic microwave background (CMB) radiation, and the effect of galactic outflows expelling large amounts of gas (Vallini et al. 2013, 2015; Pallottini et al. 2017a, 2017b).

In this letter, we reanalyze archival ALMA [C II] observations of Himiko, one of the most famous LAEs, at $z_{\text{Ly}\alpha} \sim 6.595$ (Ouchi et al. 2009, 2013; Zabl et al. 2015). Hubble Space Telescope (HST) observations with WFC3 in bands J125 and H160 (rest-frame UV at $z \sim 6.595$) show that Himiko comprises of three bright clumps (Figure 1). Its Ly$\alpha$ nebula extends over 17 kpc embedding all of the tree clumps, but its peak is located between clumps A and B (Ouchi et al. 2013). The first ALMA observations aimed at detecting [C II] emission in the Himiko system were obtained in Cycle 0 and presented in Ouchi et al. (2013), who reported a non-detection for this galaxy, placing an upper limit for the [C II] luminosity at $L_{\text{C II}} < 5.4 \times 10^7 L_\odot$. Himiko was also observed with ALMA in Cycle 1, but the analysis of these new observations have not been published so far.

In this paper we present an analysis of both Cycle 0 and Cycle 1 data. When combined together, we show that such data present a clear detection of [C II] with a luminosity expected by the local [C II]–SFR relation, and also present a clumpy morphology relative to the UV distribution.

Throughout this paper we adopt the standard cosmological parameters $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_M = 0.30$, $\Omega_\Lambda = 0.70$. The magnitude is given in the AB system.

2. Observation and Data Reduction

Himiko was observed with ALMA in band 6 in 2012 July and 2015 July as part of the programs #2011.00011.S (Cycle 0) and #2012.10033.S (Cycle 1). The former program was carried out using a semi-compact array configuration with a
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3. Results

3.1. Continuum Emission at 158 μm

The dust continuum emission at 158 μm is not detected, indicating a low dust mass content, as already discussed by Walter et al. (2012), Ouchi et al. (2013), and Ota et al. (2014). We estimate a 3σ upper limit on the continuum emission $S_{158\,\mu m} < 27 \,\mu Jy$ that results into an infrared luminosity of $L_{IR} < 3 \times 10^{10} L_{\odot}$, if assuming gray-body emission with a dust temperature of 35 K and a spectral emissivity index of 1.5 (Ota et al. 2014; Schaerer et al. 2015; Willett et al. 2015), and correcting for the effect of the CMB (da Cunha et al. 2013; Ota et al. 2014). Following the $L_{IR}$–$S_{IR}$ relation presented in Kennicutt & Evans (2012), the $L_{IR}$ translates into a 3σ upper limit on the dust obscured SFR$_{IR} < 4 \,M_{\odot}\,yr^{-1}$. We note that such upper limit on SFR$_{IR}$ is about four times lower than the total SFR$_{UV}$ inferred from the rest-frame UV emission (see Table 2), indicating that most of the star formation is unobscured by dust.

3.2. [C II] Line Emission

We detect an emission line consistent with the [C II] emission at the redshift and at the location of Himiko in the combined UV-tapered ALMA cube. Figure 1(a) shows the map of the line extracted with a spectral width of 270 km s$^{-1}$ and centered at $-185$ km s$^{-1}$ relative to the $z_{Ly\alpha}$. The peak of the [C II] emission has a significance of 5.5 σ and is located between UV clumps A and B. However, the [C II] emission is spatially resolved, hence the significance of the total emission is higher than that of the peak flux in the map and specifically, as we will see below, the total significance is $\sim 8\sigma$. The distribution of the [C II] emission, estimated using CASA’s IMFIT task, has a beam-deconvolved size of $(0.7 \pm 0.2)'' \times (0.3 \pm 0.2)''$ embedding both UV clumps A and B.

The reliability of the [C II] detection is also supported by the fact that the line is marginally detected at the same location in

maximum baseline length of 440 m. An extended-array configuration was used instead for Cycle 1 observations, with baseline lengths up to 1600 m. Both observations were performed in frequency division mode with a total bandwidth of 7.5 GHz and a spectral resolution of 0.488 kHz ($\sim$0.5 km s$^{-1}$).

The two data sets have been calibrated using the standard scripts for ALMA data reduction. We used the software CASA (McMullin et al. 2007) version 3.4.0 and 4.7.2, respectively, for the two data sets. For each data set we have generated the continuum map and cubes using the CASA task CLEAN with a natural weighting. The resulting sensitivities and angular resolution are listed in Table 1. We note that the sensitivities and angular resolution achieved in Cycle 0 images are in agreement with those presented by Ouchi et al. (2013).

We have also combined the observations of both programs to obtain maps and cubes with higher sensitivity. The imaging of the combined data set has been performed with a natural weighting, achieving a continuum sensitivity of $9 \,\mu Jy\,beam^{-1}$ and an angular resolution of $\sim 0''30$. To optimize the signal-to-noise ratio (S/N) of the [C II] detection, we have produced a spatially smoothed cube using a UV tapering of $0''2$ so as to recover the spatially extended emission. We report the properties of combined and smoothed images in Table 1.

We have registered the ALMA and HST data by matching the ALMA calibrator and foreground sources to the GAIA Data Release 1 catalog (Gaia Collaboration et al. 2016).

Table 1

| Data Set       | $t_{on}$ (hr) | Antennae | Beam (μJy) | $\sigma_{cont}$ (μJy) | $\sigma_{line}$ (μJy) |
|----------------|---------------|----------|------------|-----------------------|-----------------------|
| 2011.000115.S  | 3.2           | 25       | $0''81 \times 0''57$ | 17                     | 130                   |
| 2012.100033.S  | 3.2           | 44       | $0''28 \times 0''22$ | 11                     | 98                    |
| Combined       | 6.4           | 44       | $0''30 \times 0''31$ | 9                      | 88                    |
| Combined and UV-tapered | 6.4      | 44       | $0''48 \times 0''40$ | 11                     | 101                   |

Note. (a) On-source exposure time. (b) Number of antennae. (c) Angular resolution. (d) Sensitivity at $\lambda_{rest} = 158 \,\mu m$. (e) Sensitivity in a spectral channel of 100 km s$^{-1}$.
UV and Far-IR Properties of Himiko

| Himiko | Total | Clump A | A_{C II} (Lyα peak) | Clump B (B_{C II}) | Clump C |
|--------|-------|---------|----------------------|---------------------|---------|
| R.A. [J2000] | ... | 2:17:57.61 | 2:17:57.59 | 2:17:57.57 | 2:17:57.53 |
| Decl. [J2000] | ... | −5:08:44.96 | −5:08:44.77 | −5:08:44.87 | −5:08:44.82 |
| z_{Lyα} | 6.595 | 6.595 | 6.595 | 6.595 | 6.595 |
| EW(Lyα) [Å] | 78.84 | 68.14 | >68 | 5.20 | 6.12 |
| J_{125} [μJy/beam] | 24.99 ± 0.08 | 26.54 ± 0.04 | >26.9 | 27.03 ± 0.07 | 26.43 ± 0.04 |
| H_{i153} [μJy/beam] | 24.99 ± 0.10 | 26.73 ± 0.06 | >27.3 | 27.04 ± 0.08 | 26.48 ± 0.05 |
| SFR_{C II} [M_{⊙} yr^{-1}] | 20.4 ± 1.5 | 4.9 ± 0.2 | <3.5 | 3.1 ± 0.2 | 5.4 ± 0.2 |
| S_{Lyα, jum} [μJy] | <27 | <27 | <27 | <27 | <27 |
| L_{R} [10^{48}L_{⊙}] | <3 | <3 | <3 | <3 | <3 |
| SFR_{R} [M_{⊙} yr^{-1}] | <4 | <4 | <4 | <4 | <4 |
| z_{C II} | 6.5913 ± 0.0004 | ... | 6.5906 ± 0.0008 | 6.5915 ± 0.0003 | ... |
| Δ v_{cav-[C II]} [km s^{-1}] | −145 ± 15 | ... | −175 ± 30 | −140 ± 10 | ... |
| FWHM_{C II} [km s^{-1}] | 180 ± 50 | ... | 240 ± 80 | 70 ± 20 | ... |
| S_{C II} [μJy km s^{-1}] | 108 ± 13 | <26 | 73 ± 12 | 44 ± 8 | <26 |
| L_{C II} [10^{48}L_{⊙}] | 1.2 ± 0.2 | <0.2 | 0.85 ± 0.13 | 0.50 ± 0.10 | <0.2 |
| [C II] size [kpc] | (3.9 ± 1.1) × (1.7 ± 1.1) | ... | (3.7 ± 0.6) × (2.0 ± 0.6) | <2 | ... |

Notes.

\(^{a}\) The magnitudes of A, B, and C are from Ouchi et al. (2013), while we perform 9\arcsec 68 × 9\arcsec 41 aperture photometry for A_{C II}. The latter is contaminated by the emission of clump A. The SFR_{C IV} has been estimated from J_{125} flux using Kennicutt & Evans (2012).

\(^{b}\) The L_{R} estimates include the effect of the CMB (da Cunha et al. 2013; Ota et al. 2014).

\(^{c}\) The upper limits on the [C II] emission are estimated by assuming a FWHM = 100 km s^{-1}.

The systemic systemic redshift estimated from [C II], but fully consistent once the Lyα asymmetry is taken into account.

The integrated intensity of the [C II] line is S_{C II}Δν = 108 ± 10 μJy km s^{-1}. The resulting [C II] luminosity is L_{C II} = (1.2 ± 0.2) × 10^{48} L_{⊙}, which is about two times higher than the upper limit estimated by Ouchi et al. (2013) and Ota et al. (2014). This tension can be explained by the fact that, not having a priori knowledge of the location and velocity of the [C II] emission, Ouchi et al. (2013) and Ota et al. (2014) extracted upper limits with arbitrary line widths that did not take into account fluctuations that may have been associated with a real signal. This issue also may be common for other [C II] non-detections in the literature, in the sense that some of the upper limits estimated in the past may be too low.

The measured [C II] luminosity is comparable to those measured in some other z > 6 star-forming galaxies with similar SFR_{C IV}. In particular it is akin to the L_{C II} observed in the CR7 (Matthee et al. 2017), which is a LAE with properties similar to Himiko (SFR = 44 M_{⊙} yr^{-1}). In contrast to previous claims, our detection places Himiko along the local L_{C II}–SFR given by De Looze et al. (2014), as illustrated in Figure 2.

3.3. Multicomponent System

Given the multi-clump shape of Himiko in the rest-frame UV images and the extended [C II] emission detected in the smoothed ALMA observations, we investigate the morphology of [C II] emission and its connection with the UV and Lyα counterparts.

A channel map analysis performed on the non-smoothed combined cube reveals that the extended [C II] emission is the result of two spectroscopically distinct components (hereafter A_{C II} and B_{C II}), whose emission peak at two different locations of the system. Figure 3 shows the flux maps and spectra of the two components. The latter are extracted from an
aperture as large as the beam and centered at the peak of the emission. We identify the component $A_{[C\,\text{II}]}$ with an S/N = 5 in the channel map $-100$ km s$^{-1} < v < -365$ km s$^{-1}$. The centroid of the emission is spatially offset by $\sim0''2$ toward the west relative to the UV position of clump A, but it is fully consistent with the peak of the extended Ly$\alpha$ nebula observed by Ouchi et al. (2013). We exclude that this [C II] component is directly associated with UV clump A since the [C II] emission level at the location of the UV emission is $<\sigma_2\sigma$. As we will discuss later, positional offsets between UV and FIR line emission are not so rare in primeval galaxies (Maiolino et al. 2015; Willott et al. 2015; Carniani et al. 2017b). By fitting a 2D Gaussian profile to the channel map, we find that the [C II] emission component $A_{[C\,\text{II}]}$ is spatially resolved with a beam-deconvolved size of $(0.68 \pm 0.10)'' \times (0.41 \pm 0.10)''$, which corresponds to a physical size of $\sim3.7$ kpc. The [C II] profile peaks at $-175 \pm 30$ km s$^{-1}$ and has a line width of $240 \pm 80$ km s$^{-1}$. We measure $L_{[C\,\text{II}]} = (0.85 \pm 0.13) \times 10^{8} L_{\odot}$ that corresponds to the $\sim70\%$ of the total [C II] luminosity inferred for Himiko. The other component, $B_{[C\,\text{II}]}$, peaks at $-175 \pm 30$ km s$^{-1}$ and is narrower (FWHM = $70 \pm 20$ km s$^{-1}$) than the former component. The knot of the [C II] emission is co-aligned with UV clump B, indicating that [C II] and UV emission arise from the same region. Since components $B_{[C\,\text{II}]}$ and $A_{[C\,\text{II}]}$ overlap in velocity, the flux map $B_{[C\,\text{II}]}$ shows also a tail emission associated with the $A_{[C\,\text{II}]}$ component. The core of the $B_{[C\,\text{II}]}$ component is not spatially resolved, indicating the the [C II] line is powered by a compact source with a size of $<2$ kpc. We infer a [C II] luminosity of $L_{[C\,\text{II}]} = (0.50 \pm 0.13) \times 10^{8} L_{\odot}$.

Since around the Ly$\alpha$ we do not detect any emission close to the UV clumps A and C, we search for line emission in the ALMA cube between $-1000$ and $1000$ km s$^{-1}$ relative to Ly$\alpha$. Only a putative detection is found at $\sim-500$ km s$^{-1}$ and located at the UV position of clump C. Because of the low significance (S/N = 3.2), the emission can be spurious due to noise fluctuation, hence we consider it as a non-detection. For the two UV clumps A and C, we therefore infer an upper limit on the [C II] luminosity of $L_{[C\,\text{II}]} < 0.2 \times 10^{8} L_{\odot}$, where we assume a line width of 100 km s$^{-1}$.

We note that while the [C II] emission in clump B (= $B_{[C\,\text{II}]}$) is consistent with the local $L_{[C\,\text{II}]}$–SFR relation, clumps A, C, and $A_{[C\,\text{II}]}$ are scattered outside the local relation (Figure 2), as observed in other high-$z$ galaxies (e.g., Maiolino et al. 2015). Positional offsets between UV and [C II] emission observed both in the local and distant universe may be ascribed to spatially distinct regions of the Galaxy, and of their circumgalactic environment, characterized by different physical properties (e.g., Kapala et al. 2015; Croxall et al. 2017; Carniani et al. 2017b). For instance, the low [C II] luminosity at the location of the UV regions can be interpreted as a consequence of a local low metal content, but also in terms of strong feedback ionizing or expelling gas (Vallini et al. 2015; Katz et al. 2016; Olsen et al. 2017). Spatially offset [C II] emission may also be explained in terms of dust obscuration and/or outflowing/inflowing gas. In the former scenario [C II] is excited in situ by star formation whose UV emission is heavily dust obscured. In this context we note that Ouchi et al. (2013) infer a dust attenuation for Himiko of $E(B-V) = 0.15$, which can hide a significant fraction of star-forming regions traced by the UV emission. On the other hand, the spatially offset [C II] can be associated with a satellite clump in the process of accreting, or clumps expelled by galactic outflows; in these cases the [C II] emission is excited by the UV radiation of the closest star formation region (e.g., clump A). This last scenario is also supported by the fact that Himiko reveals a triple major merger event whose extended Ly$\alpha$ nebula emission may be powered by both star formation and galactic winds (Ouchi et al. 2013; Zabl et al. 2015).

In summary, UV and [C II] emission can trace different regions that should be treated as different subcomponents of the same system. A detailed discussion of multiple subcomponents observed in $z > 5$ star-forming galaxies, as well as their offset relative to the star-forming regions traced by the UV emission, is presented in a companion paper (Carniani et al. 2017a).

4. Conclusions

We have presented the analysis of archival ALMA data targeting the [C II] emission in the famous and luminous LAE Himiko at $z = 6.595$. We have detected the FIR line with a level of significance of $9\sigma$ for the total spatially integrated emission. The measured luminosity of the line ($L_{[C\,\text{II}]} = 1.2 \times 10^{8} L_{\odot}$) is fully consistent with that of the local $L_{[C\,\text{II}]}$–SFR relation, thereby mitigating the discrepancy with the local $L_{[C\,\text{II}]}$–SFR relation claimed in previous studies. The ALMA data reveal that the [C II] profile is blueshifted by $-145$ km s$^{-1}$ relative to the peak of Ly$\alpha$, consistent with the fact that the asymmetric, blueward profile of Ly$\alpha$ is associated with absorption by the neutral IGM. The [C II] emission is spatially resolved over $\sim4$ kpc and breaks into two subcomponents. The location of the faintest [C II] component (which is spatially unresolved) is consistent with the UV emission of clump B, which is the central UV clump out of the three clumps associated with Himiko. Instead, the brightest [C II] component is spatially extended (over 4 kpc) and its centroid is coincident with the peak of the Ly$\alpha$ nebula, and is spatially offset by $\sim0''2 (=1$ kpc) relative to the nearest (and brightest) UV clump. While the [C II] luminosity of clump B is fully consistent with the local $L_{[C\,\text{II}]}$–SFR relation, all of

Figure 2. $L_{[C\,\text{II}]}$–SFR diagram. The green line shows the relations for local star-forming galaxies by De Looze et al. (2014). The location of Himiko’s total emission is shown with the large red circle. The locations of Himiko’s subclumps are shown with red squares. We also report the location of CR7 and its subclumps. The dotted circle shows the previous $L_{[C\,\text{II}]}$–SFR relation, all of which corresponds to a non-detection. For the two UV clumps A and C, we therefore infer an upper limit.
the other clumps are scattered around the local relation, both above and below.

As already discussed in other works, the offsets between \([\text{C II}]\) and UV emission, and the multi-clump nature of the \([\text{C II}]\) emission, are likely associated with various effects, such as dust obscuration (preventing the detection of UV associated with \([\text{C II}]\) emission), strong feedback (ionizing or removing gas in intense star-forming regions), minor/major mergers, and circumgalactic accreting/outflowing gas. A more extensive discussion of the multicomponent \([\text{C II}]\) properties of high-\(z\) galaxies, their morphologies, and offsets relative to the UV emission as well as several more galaxies at \(z > 5\) are presented in a companion paper (Carniani et al. 2017a).

Finally, we note that the previous works had not detected the \([\text{C II}]\) line in Himiko, giving an upper limit a factor of two lower than our detection. This is due to the difficulty of properly estimating the upper limit without prior knowledge about the location, redshift, and width of the \([\text{C II}]\) emission. This issue also may affect other past results; for example, past upper limits on the \([\text{C II}]\) emission in other high-\(z\) galaxies may in fact be too low.

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