Magnetic states stabilization in $\text{Ni}_{51}\text{Mn}_{33.4}\text{In}_{15.6}$ Heusler alloy

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Abstract: The rate-independent stabilization of magnetic states with iterations in a Heusler alloy has been studied. The direct measurement of the adiabatic temperature change, $\Delta T_{ad}$, of a $\text{Ni}_{51}\text{Mn}_{33.4}\text{In}_{15.6}$ alloy near the magnetostructural phase transition is presented. The adiabatic temperature change at a given temperature within the temperature range of the magnetostructural transition is history dependent and varies considerably with the iteration count of the field cycle. The data show the transition from the low magnetization state to the high magnetization state during low to high (L–H) temperature change direction and from high magnetization to low magnetization state during high to low (H–L) temperature change direction require several field cycles to stabilize the $\Delta T_{ad}$ measurement, similar to the accommodation phenomenon in hysteretic materials. In the mixed magnetic state inside the first-order transition, both low and high magnetization portions of the alloy exist and it varies considerably with the induced fields. This original observation emphasizes that it is incorrect to assess the performance of a magnetic refrigeration system through a single measurement, and that achieving a stable, utilizable, adiabatic temperature change requires several field-induced transitions.

Subjects: Applied Physics; Condensed Matter Physics; Materials Science; Metals & Alloys; Physical Sciences

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PUBLIC INTEREST STATEMENT
In the advancement of alternative energy sources, magnetocaloric effect shows a promising technology for development of compact and energy efficient magnetic refrigerators. In the past 20 years, there has been a surge in research on the magnetocaloric response of materials, due mainly to the possibility of applying this effect for magnetic refrigeration close to room temperature. However, the magnetic materials available and studied by the scientific community do not yet have the necessary characteristics to be used in large scale, due to technological and/or economic restrictions. The purpose of this paper is to investigate the history-dependent behavior of the spin-reorientations and its effect on the directly measured adiabatic temperature change ($\Delta T_{ad}$) in a $\text{Ni}_{51}\text{Mn}_{33.4}\text{In}_{15.6}$ Heusler alloy.
1. Introduction

Heusler alloys have been studied in recent years due to their wide useful applications such as shape memory alloys (Krumme et al., 2015). Ni-Mn-based Heusler alloys have been the subjects of many studies because of their unusual and complex magnetic properties as well as their potential use as magnetocaloric devices (Du et al., 2007; Hütten et al., 2006; Krenke et al., 2007; Seyoum et al., 2013; Sharma, Chattopadhyay, & Roy, 2007). They are known to exhibit mixed magnetic states during the magnetostructural transformation resulting in a disorder and competing interactions between different magnetic sublattices (Aksoy, Acet, Deen, Mañosa, & Planes, 2009; Bennett et al., 2012; Ye et al., 2010). Spin reorientation mainly affects the orientation of the spins, while the magnetization magnitude remains unchanged. Such weak-order changes could be observed by few techniques and methods: neutron diffraction, X-ray diffraction, and AC resistometry (Gotoh, Ohashi, Kanomata, & Yamaguchi, 1995; Khovailo, Abe, & Takagi, 2003; Skanthakumar, Zhang, Clinton, Li, & Lynn, 1989).

Alternatively, additional peaks and anomalies in the direct measurements of the adiabatic temperature change reflect the existence of spin reorientations and the temperature range where they happen. Adiabatic temperature change is defined as the heating or cooling of magnetic material due to varying magnetic field under adiabatic conditions. Previous studies of spin-reorientations focused mainly on driving the change through controlling the alloy composition or the way it was prepared (Han et al., 2007; Krenke et al., 2006). The purpose of this paper is to investigate the history-dependent behavior of the spin reorientations and its effect on the directly measured adiabatic temperature change ($\Delta T_{ad}$) in a Ni$_{51}$Mn$_{33.4}$In$_{15.6}$ Heusler alloy.

2. Measurements and discussion

An ingot with a nominal composition of Ni$_{51}$Mn$_{33.4}$In$_{15.6}$ was prepared by induction melting pure metals in an argon atmosphere. A rectangular prism specimen, with dimensions 25 mm × 11 mm × 10 mm, was cut from the ingot. The specimen was heat treated in an argon atmosphere at 900°C for 8 h, aged at 500°C for 48 h, and then quenched in water. A photography of the sample is shown in Figure 1.

A photo of the experimental apparatus is shown in Figure 2. The experimental setup, its accuracy, and more details about the measurement procedure can be found in Ghahremani, Jin, et al. (2012), Ghahremani, Seyoum, ElBidweihy, Della Torre, and Bennett (2012). The magnetization, $M$, vs. the temperature, $T$, measurements of Ni$_{51}$Mn$_{33.4}$In$_{15.6}$ under a field as low as 100 Oe showed that the temperature range for the magnetostructural transition is 280–310 K and 290–310 K for the L–H and H–L processes, respectively (Seyoum et al., 2013).

Ghahremani et al. (2013) studied the latent heat of phase transmission in Heusler alloy. They measured an independent adiabatic temperature change ($\Delta T_{ad}$) of Ni$_{51}$Mn$_{33.4}$In$_{15.6}$ during the magnetization for L–H and H–L temperature change direction processes for a 2 T applied magnetic field as shown in Figure 3. The measurement procedure that was followed in Ghahremani et al. (2013) work was done through the following process:
(1) The sample temperature is set to an initial temperature that is 281 K for L–H process and 305 K for H–L process.

(2) The magnetic field of 2 T is applied to the sample.

(3) \( \Delta T_{\text{ad}} \) is then measured and recorded.

(4) The magnetic field is removed. The system is allowed to equilibrate. A waiting time of 20 min between two successive temperature measurements is used. A waiting time up to 130 min shows no additional differences and thus guarantees that complete equilibrium had been reached and displays the independency of each data point.

(5) The sample’s initial temperature is increased or decreased by 0.5 K for L–H or H–L processes, respectively.

(6) Repeat Steps (2–5) until reaching up to 305 K for the L–H process and 281 K for the H–L process.

This experimental procedure guarantees that each data point depends only on whether the system is transforming from ferrimagnetic to ferromagnetic to paramagnetic or vice versa. Otherwise stated, each data point depends on whether the temperature change direction is H–L or L–H. Each \( \Delta T_{\text{ad}} \) measurement is recorded only upon the application of the magnetic field within a magnetization process (Step 3). As all points were measured independently, the measurement procedure cancels any contribution of magnetic hysteresis. Therefore, the hysteresis shown in Figure 3 is thermal hysteresis only.

The measurements in Figure 3 display the conventional and inverse magnetocaloric effect peaks within the temperature range for the magnetostructural transition. For the same alloy and under different magnetic field changes, additional peaks and recurring kinks in the 290–295 K temperature
range were observed (Seyoum et al., 2013). Thus, this paper describes an investigation of the behavior of spin reorientation at 292 K, a temperature lying within that temperature range.

As seen in Figure 3, at temperature 292 K (marked by dashed line), there exist two different $\Delta T_{ad}$ for L–H and H–L processes under magnetization of 2 T magnetic field. It is worth mentioning that the $\Delta T_{ad}$ at 292 K for L–H process is negative and for H–L process is positive. In order to study the history-dependent behavior of the sample in the magnetostructural transition region, the direct measurement of $\Delta T_{ad}$ for magnetization was reiterated a sufficient number of times at 292 K for both L–H and H–L processes to obtain repeatable results. The experimental procedure is as follows:

1. The sample temperature is set to an initial temperature of 292 K.
2. The sample is then magnetized by a magnetic field of 2 T.
3. $\Delta T_{ad}$ is then measured and recorded.
4. The magnetic field is removed. The system is allowed to equilibrate.
5. For L–H process, the sample is cooled down to 280 K, and then the sample’s temperature is increased to 292 K. For H–L process, the sample is heated up to 305 K, and then decreased again to 292 K. (This step guarantees independent measurements, and removes the existence of any magnetic hysteresis.)
6. Repeat Steps (2–5) until a stable measurement is recorded.

The magnetization adiabatic temperature change ($\Delta T_{ad}$) vs. the number of iterations for the L–H and H–L temperature change directions is shown in Figures 4 and 5, respectively. For the L–H process, the measured magnetization $\Delta T_{ad}$ at 292 K for the first iteration is $-2$ K and it eventually stabilizes at $-1.57$ K. For the H–L process, the measured magnetization $\Delta T_{ad}$ at 292 K for the first iteration is $0.96$ K and it eventually stabilizes at $0.86$ K. It was found that, for this sample, nine iterations were sufficient to obtain a repeatable $\Delta T_{ad}$ measurement at 292 K.

The effect shown in Figures 4 and 5 is similar to the accommodation phenomenon in hysteretic materials (Bennett, Vajda, Atzmony, & Swartzendruber, 1996; Della Torre, 1994; Della Torre & Jin, 2009). Accommodation (reptation) is the cyclical change of minor hysteresis loops towards an equilibrium limit cycle. The effect exhibited by the $\Delta T_{ad}$ measurements in this paper is rate independent, similar to accommodation. The stabilization occurs because the spin reorientation between the mixed magnetic states within the magnetostructural transition region is asymmetrical and requires several field-induced transitions to be repeatable. Unlike the magnitude of the magnetization, the adiabatic temperature change is sensitive to the orientations of the spins. Depending on the strength

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**Figure 4.** The adiabatic temperature change, $\Delta T_{ad}$ vs. the iteration count measured during the L–H magnetization process for a magnetic field change of 2 T.
of the field change, spins will reorient to a new favorable magnetization axis. When the field is removed, not all the spins will revert back to the original orientation. After several field cycles, all the retained spins contribute to the reorientation process and the effect is now deterministic leading to a stable measurement. It is concluded that $\Delta T_{\text{ad}}$ will display a rate-independent transient-like behavior for the first few measurements before a consistent, stable measurement is obtained.

This stabilization effect, revealed for the first time in this paper, contributes to a significant difference in calculating the coefficient of performance (COP), a coefficient used to measure the performance of magnetic refrigeration systems. COP is equal to the ratio of heating or cooling provided ($Q$) to the energy dissipated ($W$). The amount of heat generated in each magnetization cycle is given by:

$$Q = m \times c_p \times \Delta T_{\text{ad}},$$

where $m$ is the mass of the sample and $c_p$ is the specific heat capacity at a given temperature. The relationship between the COP and the adiabatic temperature change is, therefore, given by:

$$\text{COP} = \frac{Q}{W} = \frac{m \times c_p \times \Delta T_{\text{ad}}}{W}.$$  

As seen in Figures 4 and 5, the calculated COP at 292 K of the 1st iteration is 27% larger than that of the 9th iteration for the L–H process, and 12% larger for the H–L process. Considering that researchers usually record $\Delta T_{\text{ad}}$ after the first iteration, and that magnetic refrigeration is a cyclical process, our findings emphasize the error involved in measuring $\Delta T_{\text{ad}}$ or calculating COP after only one field iteration/cycle. This error leads to an overestimate of the performance of a magnetic refrigeration system.

3. Conclusion

Cyclical rate-independent measurements of the adiabatic temperature change, $\Delta T_{\text{ad}}$, of Ni$_{51}$Mn$_{33.4}$In$_{15.6}$ revealed that spin reorientations within the temperature range of the magnetostructural transition (292 K) are asymmetrical, vary with the iteration count, and change step by step into an equilibrium limit cycle. These weak-order changes require several field cycles to drive the measured $\Delta T_{\text{ad}}$ into a stable measurement, similar to the accommodated minor loops of hysteretic materials. The measurements presented emphasize the overestimate of the performance of a magnetic refrigeration system using a single cycle of $\Delta T_{\text{ad}}$ direct measurement.

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