The Two-Photon Reversible Reaction of the Bistable Jumping Spider Rhodopsin-1

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ABSTRACT Bistable opsins are photopigments expressed in both invertebrates and vertebrates. These light-sensitive G-protein-coupled receptors undergo a reversible reaction upon illumination. A first photon initiates the cis to trans isomerization of the retinal chromophore—attached to the protein through a protonated Schiff base—and a series of transition states that eventually results in the formation of the thermally stable and active Meta state. Excitation by a second photon reverts this process to recover the original ground state. On the other hand, monostable opsins (e.g., bovine rhodopsin) lose their chromophore during the decay of the Meta II state (i.e., they bleach). Spectroscopic studies on the molecular details of the two-photon cycle in bistable opsins are limited. Here, we describe the successful expression and purification of recombinant rhodopsin-1 from the jumping spider Hasarius adansoni (JSR1). In its natural configuration, spectroscopic characterization of JSR1 is hampered by the similar absorption spectra in the visible spectrum of the inactive and active states. We solved this issue by separating their absorption spectra by replacing the endogenous 11-cis retinal chromophore with the blue-shifted 9-cis JSiR1. With this system, we used time-resolved ultraviolet-visible spectroscopy after pulsed laser excitation to obtain kinetic details of the rise and decay of the photocycle intermediates. We also used resonance Raman spectroscopy to elucidate structural changes of the retinal chromophore upon illumination. Our data clearly indicate that the protonated Schiff base is stable throughout the entire photoreaction. We additionally show that the accompanying conformational changes in the protein are different from those of monostable rhodopsin, as recorded by light-induced FTIR difference spectroscopy. Thus, we envisage JSR1 as becoming a model system for future studies on the reaction mechanisms of bistable opsins, e.g., by time-resolved x-ray crystallography.

INTRODUCTION

Animal opsins are light-sensitive G-protein-coupled receptors (GPCRs) mainly involved in vision and circadian clock entrainment (1). In these GPCRs, photon absorption results in the isomerization of a retinal chromophore covalently linked to the protein through a protonated Schiff base. Retinal isomerization leads to structural changes in the receptor, resulting in the recruitment and activation of G-proteins and other downstream signaling cascades.

According to the stability of the photoactivated state, opsins can be classified as monostable or bistable. In monostable opsins, deprotonation and subsequent hydrolysis of the Schiff base during Meta II decay leads to the eventual loss of the retinal chromophore (bleaching). Once released, the all-trans retinal is reisomerized in the retinal pigment epithelium and reconstituted into an opsin obtaining a functional rhodopsin (2). In contrast, in bistable opsins—such as jumping spider rhodopsin-1 (JSR1) (3)—retinal remains in the protein binding pocket throughout the entire photoreaction (4), and its isomerization leads to the formation of a thermally stable active state (acid-Meta) (4–7), suggesting the presence of a protonated Schiff base (4,6). Illumination of this state recovers the original inactive ground state (Rho). Hence, these bistable opsins exhibit both forward and backward photoreactions, i.e., the retinal isomerizes back and forth between cis and trans configurations upon repeated illumination (7).

The fast rate of these photoreactions hinders our ability to understand the activation mechanism of bistable opsins at the molecular level. Previous studies on squid and octopus rhodopsins revealed numerous intermediates in the forward...
reaction of the cycle (all-trans half-cycle: Rho > Batho > Lumi > Meso > acid-Meta), and an additional state between Meso and Meta (t-Meta) has been proposed for octopus rhodopsin (8–10). However, insights about the backward reaction—11-cis half-cycle—are scarce. In octopus rhodopsin, this reaction has been suggested to consist of only two detectable intermediates (acid-Meta > I1 > I2 > Rho) and to be considerably slower than the forward reaction (11). In the case of squid rhodopsin, it has been suggested that the Schiff base de- and reprotonates during the recovery of the ground state (12).

There are several issues that complicate spectroscopic studies of bistable opsins, such as being able to disentangle the similar absorption spectra of the ground and active states or the absence of a model system—like bovine rhodopsin for monostable opsins (2)—that can be recombinantly expressed and stably purified to yield large quantities of functional protein (13). To date, bistable opsins for biophysical studies are sourced from native retinae, precluding protein engineering and detailed investigations (14–17).

Here, we report the biochemical and biophysical characterization of recombinantly produced JSR1, a potential model for bistable opsins. JSR1 is a green-light sensitive visual pigment found in the rhabdomeric photoreceptor cells of Hastrarius adansoni (3). The endogenous chromophore of JSR1 is the 11-cis retinal isomer, with an absorbance maximum at 535 nm in the ground as well as in the all-trans photoprotein Meta state (3). To improve the spectroscopic separation between inactive and active states of the protein, we used the 9-cis retinal analog because of its hypochromic shift (isorhodopsin, JSiR1). Furthermore, we have characterized the forward and backward photoreactions of JSiR1 under ambient conditions using time-resolved ultraviolet-visible (UV-Vis) spectroscopy, identifying different intermediates and their time constants of rise and decay. Additionally, vibrational spectroscopy provided evidence for structural changes of JSiR1 upon photoactivation on the level of single bonds.

This detailed characterization of the photoactivation process of JSR1 sheds light on the mechanistic basis of bistability in light-sensitive GPCRs. We also expect that these results will assist in establishing JSR1 as a model system for the study of novel light-controlled molecular switches in optogenetic applications.

METHODS

Cloning and expression

The cDNA of JSR1 (18) from Hastrarius adansoni (AB251846) was cloned into a pcDNA 3.1(+) vector (Thermo Fisher Scientific, Carlsbad, CA) and stably expressed in HEK 293 GnTI—(gift from Khorana (19)) in Dulbecco’s modified Eagle’s medium (Thermo Fisher Scientific) supplemented with 10% fetal bovine serum and G418 (Thermo Fisher Scientific) as a selection marker. The stable clones were adapted to suspension systems and cultured in FreeStyle 293 Expression Medium (Thermo Fisher Scientific), reaching a final cell count of 2.5–3 × 10^6 cell/mL. The cultures were then harvested by centrifugation at 10,600 × g for 10 min at 4°C.

Membrane preparation and protein purification

The centrifuged cells were resuspended in 50 mM Hepes, 140 mM NaCl, and 3 mM MgCl2 with protease inhibitors (Roche Complete EDTA-free Protease Inhibitor Cocktail; Sigma-Aldrich, Buchs, Switzerland) and mechanically homogenized. The membrane suspension was reconstituted with 9-cis retinal (98% grade; Sigma-Aldrich) or 11-cis retinal (98% grade, from National Eye Institute, National Institutes of Health, Bethesda, MD) for 16 h at 4°C, after which it was solubilized with n-dodecyl-β-D-maltopyranoside (DDM Anagrade; Anatrace, Maumee, OH) for 2 h at 4°C. The solubilized membranes were centrifuged at 100,000 × g for 1 h at 4°C. The supernatant was then supplemented with cyanogen bromide-activated Sephrose 4B resin (GE Healthcare Life Science, Freiburg, Germany) bound with anti-1D4 antibody (Cell Essentials, Boston, MA). The resultant slurry was loaded onto a BioRad Glass Econo-Column (BioRad, Hercules, CA) and washed off with 50 mM Hepes, 140 mM NaCl, and 0.195 mM DDM (opsin buffer). Protein elution was carried out overnight at 4°C in opsin buffer with 800 μM 1D4 elution peptide (Peptide 2.0, Chantilly, VA). The eluted receptor was concentrated using an Amicon Ultra 100 kDa (MilliporeSigma, Burlington, MA) concentrator in a bench top centrifuge at 1083 × g for 1 h at 4°C. The concentrated protein was loaded onto a Superdex 200 Increase 10/300 GL (GE Healthcare Life Science) SEC column using an FPLC system (Äkta Explorer 10; GE Healthcare Life Science). The main fractions containing the receptor were pooled and concentrated. The sample preparation was further assessed by sodium dodecyl sulphate-polyacrylamide gel electrophoresis (SDS-PAGE) and UV-Vis spectroscopy. The purification and the other experiments have been performed under 640 nm light.

UV-Vis spectroscopy, illumination, and acid denaturation assay

A commercially available spectrometer (UV-2401PC; Shimadzu, Kyoto, Japan) was used to record spectra of JSiR1 samples at 20°C. For all the experiments reported here, purified JSiR1 was used. Purity is assed with the optical density (OD) 280/505 nm (εmax) ratio, where 2.5–2.8 was observed. SDS-PAGE gels and size-exclusion chromatogram traces were used to evaluate sample quality. Protein concentration was calculated by the Beer-Lambert law, where εJSIR1 is ~32,660 M⁻¹ cm⁻¹ (see below), and the molecular weight was calculated from the protein sequence as 44 kDa. Illumination was carried out in the dark with 495 nm (low-pass filter) on a projector lamp at maximal intensity for 10 min (lamp parameters 150 W, 2 A, 50/60 Hz, 220–240 V, irradiance 5 W/cm²). In the acid denaturation assay, purified JSiR1 at ~OD 0.05 was supplemented with 4.275 M HCl, and a spectrum was recorded soon after (20).

Extinction coefficient of 9-cis retinal oxime

The extinction coefficient of the 9-cis retinal oxime was determined in the opsin buffer as described. A standard curve was calculated with different concentration of 9-cis retinal oxime. 1 M hydroxylamine was dissolved in the opsin buffer (opsin buffer H). This solution was also used to baseline the spectrophotometer. A stock of 4 μM 9-cis retinal was created in the opsin buffer and incubated for 20 min in ice to complete the oxime formation. From this stock, 10 dilutions of 9-cis retinal oxime were created in opsin buffer H (100, 50, 25, 12.5, 6.25, 3.12, 1.56, 0.78, 0.39, 0.19 μM), and their absorbance was measured. Collected data were interpolated where the slope of the line equals the molar absorptivity coefficient (ε × l, where l is the pathlength), and therefore ε = slope/l (Fig. S5).
Extinction coefficient of JSiR1 and hydroxylamine assay

JSiR1 in opsin buffer was treated with 300 mM hydroxylamine at final 0.119 OD. The sample with hydroxylamine was measured at different time points for 3 h at 20°C.

Δ360 is defined as the difference in absorbance of “hydroxylamine-treated sample—ground state” (how much retinal is released), and Δ505 is the difference in absorbance of “ground state—hydroxylamine-treated sample” (amount of opsin bleached). Assuming that one molecule of JSiR1 releases one molecule of 9-cis retinal, Δ360/Δ505 = εretinal/εJSiR1, and therefore εJSiR1 = (Δ505 × εretinal)/Δ360, where εretinal is 46,930 M⁻¹ cm⁻¹, as calculated above (Fig. S3). Δ360 and Δ505 are plotted against time, showing a simultaneous increase (or release) of 9-cis retinal, respectively. One-phase association nonlinear regression was used for Δ360, whereas one-phase decay was used for Δ505 (Fig. S3 B). The linear regression between Δ360/Δ505 calculates the slope of the intercept (Fig. S3 C). This method is based on (21,22).

Extinction coefficient of JSR1 photoproduct

Upon illumination, JSIR converts into a stable and heterogeneous photoproduct (JSR1 photoproduct), as shown by HPLC analysis (see Results). Therefore, ODground /εJSiR1 = ODphotoproduct /εphotoproduct, and so εphotoproduct = (ODphotoproduct × εground state) /ODground state. To validate this reasoning, the ratio ODphotoproduct/ODground state = εphotoproduct/εground state has been calculated and plotted through different sample concentrations (4.0, 2.4, 2.0, 1.4, and 1.0 μM) (Fig. S5). In addition, the slope of intercept corresponds to ODphotoproduct/ODground state (Fig. S5).

HPLC

Extraction of the chromophore from non- and irradiated JSiR1 was based on Groenendijk et al. (23) and analyzed by high-performance liquid chromatography (HPLC, Hewlett Packard Series 1050, diode array detector; Hewlett Packard, Palo Alto, CA). About 100 μL of purified JSiR1 at OD 0.2 was treated with 250 μL 20% cold methanol and then with 50 μL 1 M hydroxylamine (NH₂OH) (pH 7.0) and incubated on ice for 10 min. Finally, n-hexane (400 μL) was added to separate retinal oximes from the rest. The sample was aliquoted in dark glass vials, injected onto a YMC-Pack silica column (particle size 5 μm, volume 150 × 60.0 mm), and eluted with n-hexane containing 15% ethyl acetate and 0.15% ethanol at a flow rate of 1.0 mL/min. Data were recorded by a diode array detector at 360 nm. The results were validated through comparison with literature and internal calibration curves of diverse retinal oximes. The area under the peaks was calculated with baseline correction in Origin (OriginLab, Northampton, MA). Area under the peak was used to compute the relative oxime extinction coefficient (24) to obtain a correct value. All the plots were created with GraphPad Prism, and MarvinView was used to draw the retinal isomers.

Time-resolved UV-Vis spectroscopy

Time-resolved UV-Vis experiments on JSR1 reconstituted with 9-cis retinal (JSR1, isorhodopsin) solubilized in 0.195 mM DDM, 50 mM HEPES, 140 mM NaCl (pH 6.5) were performed using a commercial flash photolysis setup (LKS70; Applied Photophysics, Leatherhead, UK) essentially as described in (25). Briefly, the photoconversion was induced by a short laser pulse emitted by a neodymium yttrium aluminum garnet laser (Quanta-Ray: Spectra-Physics, Santa Clara, CA), which drives an optical parametric oscillator. The emission (10 ns pulse width) was set to a wavelength of 465 nm and an energy density 3–4 mJ/cm² at the sample, leading to a photoconversion of ~6–9% of JSR1 (as determined by UV-Vis spectroscopy; data not shown). Single-shot absorption changes were recorded in a range from 400 to 620 nm in 20 nm steps (emitting 460 nm due to scattering by the photolyzing laser) on two timescales, with the faster time range (<1 ms) recorded with the light source (Xe arc lamp) in pulsed mode and the slower timescale (>10 μs) recorded with the lamp running in continuous operation. The data were merged to cover the timescale from 50 ns to 1 s. Because the exciting laser flash reversibly photoconverts a fraction of the initial proteins with 9-cis retinal to the all-trans Meta state, data were recorded in both wavelength directions, i.e., from 400 to 620 nm and vice versa. A fresh sample was used for each direction to obtain meaningful data of the Iso-to-Meta transition. After 11 laser flashes (requirement for one wavelength scan), the accumulation of Meta raised to ~66%, which was proportionally subtracted from the measured data. The resulting two data sets were intensity corrected and merged with weighting factors for each wavelength, taking account of the varying contribution of the Iso-to-Meta transition in the data set. All of this was achieved by using a nonlinear least squares regression, expecting the data set to finish with a difference spectrum of Meta-minus-Iso (Fig. S3 A).

For experiments on the Meta state, the sample was illuminated by a 473 nm diode-pumped solid-state laser (30 mW; CNC, Changchun, China) for 5 min before the experiment. All measurements were performed at 20°C using a circulating water bath (F25; JULABO, Seelbach, Germany). For the global fit of the illuminated sample, we used a model of two parallel photoreactions, one representing the forward reaction with fixed time constants (Fig. 4 B, bottom panel, gray lines) and the other with indefinite time constants for the backward reaction. All calculations were done in MATLAB (The MathWorks, Natick, MA).

FTIR spectroscopy

A highly concentrated JSiR1 at pH 6.5 sample was dried on a BaF₂ window and sealed by a second window. The protein film was rehydrated via the saturated vapor phase of a glycerol/water mixture (5:5 wt/wt). Light-induced infrared (IR) difference spectra were recorded in transmission mode using a commercial Fourier Transform infrared (FTIR) spectrometer (Vertex 80v; Bruker, Billerica, MA) as described in (25), running at a spectral resolution of 2 cm⁻¹. A spectrum was taken while keeping the sample in the dark as a background. Subsequently, spectra were recorded after 5 min of illumination by a light-emitting diode (LED) emitting at 470 nm.

Resonance Raman spectroscopy

Resonance Raman spectroscopy was performed on a LabRAM spectrometer (JobinYvon/HORIBA, Kyoto, Japan) as described in (26) but using frequency-stabilized diode-pumped solid-state lasers (CNI) that emit at 457 or 532 nm, respectively. Before the Raman recordings with the rotational cuvette, 200 μL of concentrated JSiR1 at pH 6.5 was illuminated by a 470 nm LED. For deuterium exchange, the sample was rigorously washed with D₂O containing 50 mM HEPES, 140 mM NaCl, and 0.195 mM DDM. Low-temperature experiments were done in a cryostage (TMHSS500; Linkam, Tadworth, UK). Concentrated JSiR1 was dried on a silicon crucible, rehydrated with 4 μL of an H₂O/glycerol or D₂O/glycerol mixture via the vapor phase and sealed with a glass coverslip using vacuum paste. The sample was carefully protected from light until the desired temperature was reached.

Photostationary state simulation

The population of a photocycle intermediate depends on the photoreaction upon absorption of a photon by the parent state (photolysis rate kph, its thermal decay rate kth, and on the concentration of its precursor. As examples for Rho,
\[
\frac{d[Rho]}{dt} = -k_{ph,Rho} \times [Rho] \times t + k_{ph,M2} \times [M2] \times t,
\]

where the photolysis rate \( k_{ph,s} \) depends on the quantum efficiency \( \Phi_s \), the extinction coefficient \( \epsilon_s \), the frequency \( v \), and the irradiance of the incoming light in W/cm\(^2\) (27). An irradiance of 5 W/cm\(^2\) at 495 nm was chosen to simulate the conditions used in the retinal extraction experiments. Extinction coefficients were determined using the absolute spectra displayed as Fig. S8. The system of ordinary differential equations (ODEs) was solved with a fifth-order Runge-Kutta algorithm written in Python. We neglected in our model a possible photolysis of the thermal decaying states and the Batho intermediate, permitting a bigger step size, namely 250 ns for time constants below 60 \( \mu \)s for the back reaction and 20 \( \mu \)s for time constants above 300 \( \mu \)s. Simulations were run until the equilibrium was reached (convergence). As a measure for accuracy, the mean absolute percentage error was used:

\[
\text{accuracy} = \frac{100}{3} \times \left( \frac{[\text{Meta}^{\text{HPLC}}] - [\text{Meta}^{\text{Sim}}]}{[\text{Meta}^{\text{HPLC}}]} + \frac{[\text{Rho}^{\text{HPLC}}] - [\text{Rho}^{\text{Sim}}]}{[\text{Rho}^{\text{HPLC}}]} + \frac{[\text{Iso}^{\text{HPLC}}] - [\text{Iso}^{\text{Sim}}]}{[\text{Iso}^{\text{HPLC}}]} \right).
\]

**RESULTS**

**Large-scale recombinant expression and purification**

Wild-type JSR1, tagged at the C-terminus with a 1D4 epitope (3), was stably expressed in suspension cultures of HEK293 GntTI– cells (19). During membrane preparation, JSR1 was reconstituted with 9-\( \text{cis} \) retinal (JSiR1, isorhodopsin) and purified using the 1D4 affinity system followed by SEC (see Methods). Although we also succeeded in reconstituting JSR1 with 11-\( \text{cis} \) retinal (Fig. S1 A), we used the 9-\( \text{cis} \) isomer because it leads to the formation of the blue-shifted state isorhodopsin (JSiR1), resulting in the spectral separation of the ground and photoproduct states (Fig. 1). Successful incorporation of the chromophore, sample purity, and concentration were verified by UV-Vis spectroscopy (Figs. 1 and S1). Spectroscopically active and pure JSiR1 elutes as a monodisperse and homogenous sample (Fig. S2) when subjected to size-exclusion chromatography (SEC). Typical yields are about \(~0.3 \text{ mg per gram of cell mass.}\)

**Features of JSiR1**

JSR1 reconstituted with 9-\( \text{cis} \) retinal (JSiR1) exhibits a maximal absorbance in the visible at 505 nm (isorhodopsin, Fig. 1). Irradiation with green light results in a stable hyper- and bathochromic shift of the opsin to an absorption maximum at 535 nm (3). When using 11-\( \text{cis} \) retinal, the ground and photoproduct states have both maximal absorbance at 535 nm (Fig. S1 A), as previously shown by Nagata et al. (3). Both the 9-\( \text{cis} \) and 11-\( \text{cis} \) retinals show an hyperchromic shift upon illumination, confirming the bistable nature of the opsin (4). The presence of a Schiff base between retinal and the opsin (Lys 321, JSR1 numbering) was demonstrated using an acid denaturation assay (5). In this assay, opsin molecules are denatured, and Schiff bases are exposed to solvent, showing absorption maxima at 440 nm (Fig. S1 B). To exclude unspecific Schiff bases on possible solvent-exposed lysines, JSiR1 was treated with hydroxylamine (\( \text{NH}_2\text{OH} \)) (21). When hydroxylamine cleaves the Schiff base between the retinal and Lys 321, the absorption maxima at 505 nm drops (opsin is formed), and the relative retinal oximes are formed, absorbing around 360 nm (Fig. S3 A).

The molar absorption and retinal composition of both JSR1 ground state and JSR1 photoprodut were then investigated. The extinction coefficient of JSR1 ground state was calculated assuming that one molecule of JSiR1 releases one molecule of 9-\( \text{cis} \) retinal, and thus \( \Delta 360/\Delta 505 = \epsilon_{\text{retinal}}/\epsilon_{\text{JSiR1}} \) (see Methods; (21,22)). JSiR1 ground-state extinction coefficient was found to be 32,660 M\(^{-1}\) cm\(^{-1}\) (Fig. S4). Fig. S3 C shows the correlation between increase/release of 9-\( \text{cis} \) retinal, respectively. The slope of the intercept corresponds to the previous relationship, and its negative value \((-1.37 \pm 0.05)\) confirms that when \( \Delta 360 \) decreases, \( \Delta 505 \) increases and vice versa.

As for other bistable opsins, the photoprodut of JSR1 is thermally stable and also results from equilibrium of diverse species, depending on parameters such as pH, temperature, and illumination. For the scope of this study, we consider the JSR1 photoprodut (acid-Meta or Meta) as the illuminated form of the receptor at pH 6.5 as described in Methods (6,28). By using the relation \( \text{OD}_{\text{JSiR1}}/\text{OD}_{\text{JSR1}} = \epsilon_{\text{photoprodut}}/\epsilon_{\text{photoprodut}} \), the extinction coefficient of
Meta is 37,560 M\(^{-1}\) cm\(^{-1}\), which is 1.15 times higher than the JSiR1 ground state and similar to the one of squid rhodopsin (22). The slope of intercept in Fig. S5 B corresponds to \(\text{OD}_{\text{photoproduct}}/\text{OD}_{\text{ground state}}\) ratio (1.13 ± 0.04), in other words, to the increment of the extinction coefficient in the photoproduct.

To further characterize chromophore composition in the recombinantly produced JSiR1, we isolated the retinal oximes both from Iso and Meta states (23) and analyzed the relative abundance of their conformations by HPLC. The data (Figs. 2 and S5) reveal an overall homogenous sample before as well as after light irradiation. The photoproduct composition is slightly more heterogeneous as found in other bistable opsins (29,30), likely resembling one of their features.

**Photoreaction dynamics: Time-resolved UV-Vis spectroscopy on JSiR1**

Photoactivation of opsins proceeds through a series of intermediates that can be traced by the spectral changes of retinal (31,32). To characterize the spectroscopic intermediates in the photoactivation of JSiR1, we probed transient absorption changes in the UV-Vis range (400–620 nm) across the time regime of 50 ns–900 ms after pulsed excitation. Photoactivation was accomplished by a 10 ns laser pulse tuned to the emission wavelength of 465 nm to induce the electronic excitation of the chromophore 9-cis retinal. This blue-shifted excitation wavelength (relative to the measured absorption maximum, Fig. 1) was chosen to minimize excitation of the first photointermediate whose short lifetime falls into the pulse width of the laser source. We separately probed the reaction from isorhodopsin to the thermally stable state (i.e., 9-cis to all-trans retinal isomerization) (Fig. 3 A) and the “equilibrium” reaction (Fig. 3 B). The latter corresponds to a preilluminated sample consisting of a mixture of the Meta state (with all-trans retinal) and the ground state (Rho with 11-cis retinal). Illumination of JSiR1 induces the formation of Meta as well as the reaction from Meta to Rho, generating a stable mixture of the latter two states with their concentration ratio being dependent on their lifetimes. Representative time traces at 580, 520, and 480 nm are shown in the top panel of Fig. 3. The data sets were globally fitted using a kinetic model based on a linear

**FIGURE 2** HPLC analysis of retinal oximes. Chromophore configurations of JSiR1 before and after light irradiation are shown. Retinal isomers have been grouped and reported in percentages. Raw data are reported in Fig. S6.
sequence of unidirectional reactions (see bottom panels of Fig. 3, A and B for the fitted concentration profiles).

A series of three sequential transitions is involved in the transition from Iso to the stable state (Fig. 3 A, middle and bottom panels). The resulting time constants obtained for each transition are $\tau_1 = 200$ ns, $\tau_2 = 80$ $\mu$s, and $\tau_3 = 3$ ms. It has been shown for bovine rhodopsin that Iso and Rho share a common Batho intermediate in the forward photoreaction and subsequently follow the same reaction path (33). On that account, we consider the Iso-to-Meta reaction to consist of the same intermediate states as the Rho-to-Meta transition. An early intermediate with an absorption difference maximum at 600 nm rises beyond the time resolution of our experiment (50 ns). We termed this early red-shifted intermediate Batho in accordance with the photoactivation cycle of other bistable opsins (octopus and squid rhodopsin (26,27)). Formation of the intermediate states is accompanied by the depletion of the initial state as detected by the negative absorption at around 550 nm at early timescale. The decay of Batho, with a time constant of 200 ns, results in an absorbance increase at around 480 nm, indicating the formation of a new early blue-shifted Lumi intermediate (blue trace in Fig. 3 A, middle panel). Although the Lumi intermediates in squid and octopus rhodopsin are red-shifted with respect to the 11-cis state Rho (8,32), it is blue-shifted in JSR1. Lumi decays with $\tau_2 = 80$ $\mu$s and leads to the formation of an intermediate with an absorption difference maximum at 520 nm called Meso (green trace in Fig. 3 A, middle panel). At the later stage, the absorption change decreases at around 520 nm, whereas a slight increase is observed at around 580 nm. At this time point, the thermally stable Meta state is formed after pulsed excitation, as inferred from the similarity to the steady-state difference spectrum of Meta-minus-Iso (Fig. S7 A). The observed subsequent changes could originate from a precursor to the final Meta state, as has been suggested for octopus rhodopsin (34). In this case, the photoreaction would be considerably slower than that of any other reaction of mono- or bistable opsins. However, it may well be that the thermally stable Meta is excited by the probing light, as observed for the long-lived P4 state of channelrhodopsin-2 (35). Omitting the changes that occur after 20 ms, the time constant for the decay of Meso is 3 ms, similar to squid rhodopsin (36).

In bistable opsins, illumination of the Meta state leads to isomerization of all-trans retinal to the 11-cis configuration corresponding to the native Rho state (11,12). We characterized this process by generating a Meta-Rho equilibrium by illumination of the initial JSR1 state with an intense laser emitting at 473 nm. The photoreaction of the Meta and Rho states was initiated by applying pulsed laser excitation. The time-resolved absorbance changes are depicted in Fig. 3 B. The photoreaction of the preilluminated sample also involves Batho (600 nm), Lumi (480 nm), and Meso (510 nm) states, reasoning for the same reaction path of Iso and Rho. It is evident that the light-induced absorption changes of the Meta-Rho mixture (Fig. 3 B) exhibit distinct differences as compared to the Iso-to-Meta reaction (Fig. 3 A). The most striking deviation is the absence of a strong positive absorption change at the end of the data recording, which indicates that the interconversion of Meta and Rho does not result in any change of absorbance. This fact reflects the similar absorption characteristics of Meta and Rho (Fig. S1). The Meta-Rho depletion, represented by the early negative absorption change, is red-shifted compared to the Iso-to-Meta reaction, representing the spectral red shift from Iso (505 nm) to Rho/Meta (535 nm). Due to this shift, Meso decreases in absorbance, and Batho is barely visible. In contrast, Lumi appears more pronounced. The rise of the time trace at 480 nm, characteristic for Lumi, clearly shows biphasic behavior that is even more pronounced in the ground-state depletion represented by the kinetic trace at 520 nm (Fig. 3 B, top). Because the laser flash excites both the Meta and the Rho states, these photoreactions are overlaid, leading to the biphasic behavior. Meso, on the other hand, forms and decays in a monoeXponential way.

To disentangle forward and reverse reactions, we searched for kinetic signatures that emerge on top of those of the activation process. Global analysis revealed that two additional transitions with time constants of 2 and 60 $\mu$s were necessary to fit the data, taking account of the biphasic rise and decay of Lumi (Fig. 3 B, bottom panel). This observation is indicative of a photoreverse reaction consisting of two intermediate states, further referred to as $M_1$ and $M_2$. Although the proposed recovery reaction of octopus rhodopsin involves two intermediates and is considerably slower than the activation process (11), the opposite is true for JSR1. It has been proposed for squid rhodopsin that the retinal Schiff base de- and reprotonates during the regeneration of Rho, indicated by an intermediate with an absorption maximum of 380 nm (12). We have not observed such a blue-shifted intermediate in the photoreaction of JSR1, which argues against transient deprotonation of the retinal Schiff base during the Meta-to-Rho reaction. This conclusion is in agreement with our resonance Raman experiments (vide infra).

Based on our time-resolved experiments, we propose a model of the two-photon cycle of JSR1 (Fig. 4). The absorption maxima of the intermediates were determined by calculating absolute spectra (Fig. S8) from the difference absorption experiments. In the next step, we validated this model by a kinetic simulation that should reproduce the relative contributions of the intermediate states as determined by retinal extraction and HPLC analysis (Fig. 2). The reactions were simulated by a system of coupled ODEs excluding photoreactions of the thermal states. These ODEs depend on the irradiance, the quantum efficiencies of the intermediate states (Iso, Rho, and Meta), and their extinction coefficients (see Methods for details). Translating
the photoconversion of 6–9% of Iso in the time-resolved experiment yields a quantum efficiency of 0.2–0.3 for the Iso-to-Meta transition, which is comparable to other mono- and bistable opsins (for the simulation, a value of 0.25 was used) (22,37) (for a detailed analysis, see Supporting Materials and Methods). An accuracy above 90% is considered to be in accordance with the HPLC data because of the experimental error in the determination of the retinal isomers. Fig. 5 shows the accuracy of the different scenarios, assuming a typical quantum efficiency of 0.7 for Rho (22,37). For time constants below 10 μs, the formation of Meta is so fast that its population depends only on the Meta-to-Rho quantum efficiency. Above 10 μs, the duration of the reverse reaction starts to take influence, decreasing drastically the overall accuracy. This indicates that Meta can only accumulate after continuous illumination if the reverse reaction is considerably faster than the photoactivation process. Best results are achieved with Meta-to-Rho quantum efficiencies of 0.4 and 0.5 and for time constants of less than 60 μs for the reverse reaction from Meta to Rho.

Structural changes associated with photoswitching: Vibrational spectroscopy on JSiR1

Resonance Raman spectroscopy is a technique that selectively enhances vibrations of the retinal. Under continuous illumination, JSiR1 forms a photostationary mixture consisting of Rho and Meta as well as other intermediate states.

To reduce their contribution in a Raman spectrum, we used a preilluminated JSiR1 solution and a rotational cuvette always providing a sample in the laser spot where all the intermediates have decayed. The top red spectrum in Fig. 6, recorded with a probing wavelength of 532 nm, shows a C=C stretch peak at 1536 cm⁻¹ that matches the correlation to the electronic absorption appearing in the UV-Vis range (Fig. 1) (38). The fingerprint region (1100–1400 cm⁻¹) exhibits bands that are mostly due to C-C stretching modes coupled to C-C-H bends and is very sensitive to the isomeric state of the retinal (33). The peak pattern at around 1201 cm⁻¹ bears great similarity to the spectrum of the (light-adapted) ground-state of bacteriorhodopsin (39), i.e., all-trans is the predominant configuration, which is characteristic for the Meta state of JSR1. However, hydrogen-out-of-plane (HOOP) modes are usually weak for all-trans Meta states, whereas they are more pronounced in Rho, reasoning for a distortion of the polyene chain in the 11-cis configuration (34,40,41). Based on this, the intense HOOP band at 967 cm⁻¹ bears great similarity to the spectrum of the (light-adapted) ground-state of bacteriorhodopsin (39), i.e., all-trans is the predominant configuration, which is characteristic for the Meta state of JSR1. However, hydrogen-out-of-plane (HOOP) modes are usually weak for all-trans Meta states, whereas they are more pronounced in Rho, reasoning for a distortion of the polyene chain in the 11-cis configuration (34,40,41). Based on this, the intense HOOP band at 967 cm⁻¹ bears great similarity to the spectrum of the (light-adapted) ground-state of bacteriorhodopsin (39), i.e., all-trans is the predominant configuration, which is characteristic for the Meta state of JSR1. However, hydrogen-out-of-plane (HOOP) modes are usually weak for all-trans Meta states, whereas they are more pronounced in Rho, reasoning for a distortion of the polyene chain in the 11-cis configuration (34,40,41). Based on this, the intense HOOP band at 967 cm⁻¹ bears great similarity to the spectrum of the (light-adapted) ground-state of bacteriorhodopsin (39), i.e., all-trans is the predominant configuration, which is characteristic for the Meta state of JSR1. However, hydrogen-out-of-plane (HOOP) modes are usually weak for all-trans Meta states, whereas they are more pronounced in Rho, reasoning for a distortion of the polyene chain in the 11-cis configuration (34,40,41). Based on this, the intense HOOP band at 967 cm⁻¹ bears great similarity to the spectrum of the (light-adapted) ground-state of bacteriorhodopsin (39), i.e., all-trans is the predominant configuration, which is characteristic for the Meta state of JSR1. However, hydrogen-out-of-plane (HOOP) modes are usually weak for all-trans Meta states, whereas they are more pronounced in Rho, reasoning for a distortion of the polyene chain in the 11-cis configuration (34,40,41). Based on this, the intense HOOP band at 967 cm⁻¹ bears great similarity to the spectrum of the (light-adapted) ground-state of bacteriorhodopsin (39), i.e., all-trans is the predominant configuration, which is characteristic for the Meta state of JSR1. However, hydrogen-out-of-plane (HOOP) modes are usually weak for all-trans Meta states, whereas they are more pronounced in Rho, reasoning for a distortion of the polyene chain in the 11-cis configuration (34,40,41).
for different states (8,42,43). On this basis, the band at 953 cm\(^{-1}\) can be assigned to isorhodopsin. Also, the fingerprint band pattern with peaks at 1144, 1217, and 1231 cm\(^{-1}\) shows similarity to a 9-cis configuration. Schiff base frequency is the same as in Meta-Rho (inset). Bottom spectrum was taken at 183 K (\(-90^\circ\) C) using a pump/probe scheme with 532/457 nm lasers, respectively. The C=C stretch vibration at 1543 cm\(^{-1}\) correlates with the visible absorption maximum of 490 nm of Lumi. Schiff base vibration is shifted to a higher frequency (inset).

By increasing the temperature, thermal transitions after Batho can take place. For octopus rhodopsin, the Lumi intermediate is stable between \(-115\) and \(-65^\circ\)C (45). The bottom blue spectrum in Fig. 7 was recorded at \(-90^\circ\)C. A Raman probe wavelength of 457 nm was chosen for selective resonance enhancement of the blue-shifted Lumi intermediate. In addition, a laser with 532 nm was used to initiate the photoconversion from Iso to Lumi. The spectrum shows a strong C=C band at 1543 cm\(^{-1}\), which matches the electronic absorption maximum of Lumi at 490 nm. Another characteristic feature of Lumi is the HOOP mode at 957 cm\(^{-1}\) (42,46). There is a small shoulder at 1617 and a distinct peak at 1650 cm\(^{-1}\) in the Schiff base region above 1600 cm\(^{-1}\). Upon \(\text{H}_2\text{O}/\text{D}_2\text{O}\) exchange, the latter band shifts to 1630 cm\(^{-1}\), indicating that the Schiff base is also protonated in Lumi (Fig. 6 inset). Interestingly, the C=N stretching frequency in \(\text{D}_2\text{O}\) is higher in Lumi than compared to the other states, arguing for an increase of the force constant of this bond. This is contrary to the behavior in octopus rhodopsin (47,48).

By analyzing the bandwidths of the C=N-H vibrational bands of the retinal Schiff base, it is possible to identify the existence of a neighboring water molecule (49). Resonance energy transfer between the Schiff base vibration and the bending mode of \(\text{H}_2\text{O}\) is broken upon exchange of \(\text{H}_2\text{O}\) for \(\text{D}_2\text{O}\). This leads to a narrowing of the Schiff base band upon H/D exchange, which can be observed for the Meta-Rho (from 15.1 \pm 1.0 cm\(^{-1}\) in \(\text{H}_2\text{O}\) to 12.3 \pm 1.8 cm\(^{-1}\) in \(\text{D}_2\text{O}\)) and Iso/Batho (from 12.4 \pm 0.3 cm\(^{-1}\) in \(\text{H}_2\text{O}\) to 9.7 \pm 2.5 cm\(^{-1}\) in \(\text{D}_2\text{O}\)) spectra in Fig. 6. This suggests a water molecule in close vicinity to the Schiff base of Iso/Batho.
base in these states. This effect is less pronounced for the Lumi state, but the analysis is more error prone (13.9 ± 0.5 cm$^{-1}$ in H$_2$O to 12.2 ± 3.7 cm$^{-1}$ in D$_2$O). Thus, we cannot unequivocally conclude on the presence of the water molecule close to the Schiff base in the Lumi state. However, the similar deuterium isotope shift in all investigated states, namely Iso, Rho, Meta, Batho, and Lumi, implies that the hydrogen bond strength of the Schiff base to the water part of the counterion system remains constant (from JSR1 structure, unpublished data).

Although resonance Raman spectroscopy focuses on the retinal chromophore, infrared spectroscopy reveals structural information of the entire protein. We performed light-induced FTIR difference spectroscopy on the 9-cis bound JSR1 to investigate the structural rearrangements in the protein backbone and protonation changes of amino acids as well as conformational changes of the retinal upon light activation (50,51). In the difference spectrum (Fig. 7, red spectrum), negative bands (−) correspond to vibrations in isorhodopsin (dark state) and positive bands (+) to the Meta-Rho mixture (illuminated state).

Some of the bands can be assigned to the retinal chromophore. The 1559(−),1541(−)/1527(+) cm$^{-1}$ feature is attributed to the C=C stretch of Iso and Meta-Rho, respectively. This assignment is aided by comparison to the differences in the Raman spectra of Meta-Rho and Isorhodopsin (Fig. S9). A negative band at 1519 cm$^{-1}$ is missing in the FTIR difference spectrum, arguing for its attribution to Batho in the Raman spectrum (Fig. 6). In the fingerprint region, the three negative bands at 1148, 1208, and 1239 cm$^{-1}$ are characteristic for the 9-cis configuration of Iso (52). The negative HOOP mode at 954 cm$^{-1}$ can be attributed to Iso, whereas the positive feature at 967 cm$^{-1}$ indicates that Rho is also formed in addition to Meta, as shown for octopus rhodopsin (42). The most intense difference bands can be found in the region between 1600 and 1700 cm$^{-1}$ containing mostly amide I (C=O stretching) vibrations. These originate from the protein backbone and thus reflect structural rearrangements. Contributions from the C=N-H vibration of the retinal Schiff base can be ruled out from the comparison with our Raman measurements (Fig. 6), which showed that the C=N-H stretching frequency is identical for the Rho, Iso, and Meta states. Thus, the corresponding vibrational bands cancel in the FTIR difference spectrum of these intermediates (8,42).

The positive bands in the FTIR spectrum, corresponding to the Meta and Rho states, do not show evident similarities to published spectra of bovine and octopus rhodopsin (42,43). This is expected because the difference spectrum reflects the vibrational changes corresponding to the 9-cis to all-trans/11-cis transition (Iso-to-Meta-Rho) for JSR1 and not to an 11-cis to all-trans (Rho-to-Meta) as for visual rhodopsin. It was found for bovine and octopus rhodopsin that Rho and Iso predominately differ in the retinal moiety and thus an infrared difference spectrum does not exhibit strong bands above 1600 cm$^{-1}$ (42,53). This suggests that the pronounced changes in amide I vibrations are linked to the Iso-to-Meta transition, meaning that these correspond to structural changes of the opsin moiety upon photocactivation.

Hydrogen-bonding changes and protonation reactions of carboxylic acid residues can be monitored in the region above 1700 cm$^{-1}$. JSR1 does not display any distinct bands in this region, except for a weak feature at 1712 cm$^{-1}$. This observation is reminiscent to bovine rhodopsin, in which Meta I does not show pronounced bands in this region. Yet, strong difference bands appear in the signaling state Meta II because of the deprotonation of the Schiff base (Fig. 7, black spectrum). A similar behavior to JSR1 was observed in the Meta state of octopus rhodopsin (42,54). Taken together, our spectroscopic results strongly suggest that the Schiff base remains protonated throughout the two-photon cycle of JSR1.

**DISCUSSION**

The inactive and active states (Rho and Meta) of JSR1 have similar absorption characteristics, making spectroscopic measurements a hard-methodical challenge. Reconstitution of JSR1 with 9-cis retinal resulted in the formation of the blue-shifted inactive isorhodopsin (JSiR1, Fig. 1), which permits the spectroscopic discrimination of ground and photoproduction states. We succeeded at producing highly homogeneous samples of JSiR1 for time-resolved experiments. Illumination of the isorhodopsin form of JSR1 results in the formation of the thermally stable Meta state within 3 ms. Transition to this state involves three intermediates, namely Batho (590 nm), Lumi (490 nm), and Meso (510 nm). The photorecovery consists of at least two intermediates and is considerably faster than the activation process, which is confirmed by photostationary state simulations. Although we have kinetic data only for the Iso-to-Meta transition, it can be concluded from the measurements on the illuminated JSR1 that Rho follows the same reaction path as Iso, as observed in other opsins (8,33). A unique feature of JSR1 is the relatively large blue shift of the Lumi intermediate. However, deprotonation of the Schiff base does not take place at this stage because such an event would shift the absorption maximum far lower, as seen in other rhodopsins (e.g., bovine rhodopsin, octopus rhodopsin, or bacteriorhodopsin (31,32,38)) and in JSR1 deprotonated mutants (unpublished data). This conclusion is solidified by resonance Raman spectroscopy on the Lumi state, in which the C=N-H vibration of the retinal Schiff base exhibits a deuterium isotope effect, implying the presence of a proton. Based on these initial experiments, we propose a photocycle model of the two-photon reaction of JSR1 as illustrated in Fig. 4.
Our photostationary state simulations show that the lower quantum efficiency in concert with a faster photorecovery ensures that Meta accumulates after illumination. Using our derived reaction scheme of JSR1 (Fig. 4), we were able to quantitatively reproduce the ratio of isomeric states determined by retinal extraction and HPLC analysis, providing additional support for our photocycle model. Assuming that the quantum efficiency of Rho in JSR1 is also 0.7, our simulations suggest that the quantum efficiency for the photorecovery is also between 0.4 and 0.5.

Light-induced FTIR difference spectroscopy revealed structural insights into the mechanism of photoactivation. The formation of Meta is accompanied by structural changes in the protein backbone, as is the case in bovine rhodopsin upon its conversion to metarhodopsin II, which is the state able to activate the G-protein (55). Spectroscopically, the change of the latter state can be identified by the deprotonation of the Schiff base, resulting in a distinct blue shift of the electronic transition from 498 to 380 nm (31). However, the Meta state of JSR1 neither exhibits such a shift nor shows carboxylic bands in the FTIR difference spectrum equivalent to metarhodopsin II (43). Furthermore, our resonance Raman experiments have shown that retinal Schiff base is protonated in the Meta state. The similar shift upon deuteration in the inactive and active states indicates that the hydrogen bond of the Schiff base proton to the water of the counterion system (observed in the crystal structure of JSiR1, unpublished data) (56) remains constant. A similar behavior is found in octopus rhodopsin, another invertebrate rhodopsin in which the Schiff base is also reported not to deprotonate (42,54). The band at 1712 cm$^{-1}$ is therefore due to a protonated carboxylic group residing in different environments in the Iso and Meta states, such as Asp147 in the DRY motif of the octopus rhodopsin (57). Our results strongly favor a scenario in which retinal isomerization induces structural changes in JSR1, eventually leading to the formation of Meta and then G-protein activation but without requiring deprotonation of the Schiff base.

We believe that our spectroscopic characterization of the two-photon cycle of JSiR1 forms the basis for future ultrafast spectroscopic and structural studies on bistable opsins. High expression yields and room temperature stability identify JSiR1 as an excellent sample for time-resolved serial crystallographic studies at synchrotrons and x-ray free electron lasers (58,59). Thus, we expect that JSR1 may serve as a model protein for further our understanding of the molecular basis of bistability in light-sensitive GPCRs and for designing new light-sensitive switches for optogenetic applications.

**SUPPORTING MATERIAL**

Supporting Material can be found online at https://doi.org/10.1016/j.bpj.2019.02.025.

**AUTHORS CONTRIBUTIONS**

G.F.X.S., J.H., and E.L. conceived the study. N.V., M.K., and E.L. cloned, expressed, and purified JSR1. They also performed the biochemical analysis of the sample and relative calculations. D.E. performed time-resolved UV-Vis, FTIR, and Raman spectroscopic experiments, photostationary state simulations, and the global fit analysis. All the authors contributed to data analysis, interpretation, and writing the manuscript.

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