REFINEMENT OF STRICHARTZ ESTIMATE FOR AIRY EQUATION IN NON-DIAGONAL CASE AND ITS APPLICATION

SATOSHI MASAKI AND JUN-ICHI SEGATA

ABSTRACT. In this paper, we give an improvement of the non-diagonal Strichartz estimate for Airy equation by using a Morrey type space. As its applications, we prove the small data scattering and existence of a special non-scattering solutions, which are minimal in suitable sense, to the mass-subcritical generalized Korteweg-de Vries (gKdV) equation. Especially, a use of the refined non-diagonal estimate removes several technical restrictions on the previous work [24] about the existence of the special non-scattering solution.

1. Introduction

In this paper, we consider space-time estimates for a solution of the Airy equation

\[
\begin{cases}
\partial_t u + \partial_x^3 u = 0, & t, x \in \mathbb{R}, \\
u(0, x) = f(x), & x \in \mathbb{R},
\end{cases}
\]

where \(f : \mathbb{R} \to \mathbb{R}\) is a given data. After the pioneering work by Strichartz [29], the space-time estimate for the dispersive equation has been studied by many authors in several directions (see for instance [7, 15] for the historical background of this topic). As for the Schrödinger equation, the Strichartz estimate for (1.1) is well-known (see [15] for instance). Grünrock [11] and the authors [23] extended the Strichartz estimate for (1.1) to the hat-Lebesgue space, more precisely we obtained the following estimate.

**Theorem 1.1** (generalized Strichartz’ estimate [11, 23]). Let \((p, q)\) be a pair satisfying either \((p, q) = (\infty, 2), (4, \infty)\) or

\[
0 \leq \frac{1}{p} < \frac{1}{4}, \quad 0 \leq \frac{1}{q} < \frac{1}{2} - \frac{1}{p}.
\]

Then, there exists a positive constant \(C\) depends only on \(\alpha\) and \(s\) such that the inequality

\[
\|\|\|

\[
(1.2) \quad \|\|\|\|
\]

holds for any \(f \in L^\alpha\), where \(\alpha\) and \(s\) are given by

\[
\frac{1}{\alpha} = \frac{2}{p} + \frac{1}{q}, \quad s = -\frac{1}{p} + \frac{2}{q}.
\]

2000 Mathematics Subject Classification. Primary 35Q53; Secondary 35B40.

Key words and phrases. generalized Korteweg-de Vries equation, scattering problem.
and the space $\hat{L}^\alpha$ is defined for $1 \leq \alpha \leq \infty$ by

$$\hat{L}^\alpha = \hat{L}^\alpha(\mathbb{R}) := \{ f \in \mathcal{S}'(\mathbb{R}) \mid \| f \|_{L^\alpha} = \| \hat{f} \|_{L^{\alpha'}} < \infty \},$$

and $\hat{f}$ stands for Fourier transform of $f$ in $x$, and $\alpha'$ denotes the Hölder conjugate of $\alpha$.

The generalized Strichartz’ estimate (1.2) is shown by interpolating the endpoint cases $(p, q) = (\infty, 2)$, $(4, \infty)$, which corresponds to the well-known Kato’s smoothing and Kenig-Ruiz estimates, and the diagonal case $p = q \in (4, \infty]$. We refer the estimate in the diagonal case to as a Stein-Tomas estimate.

The aim of this paper is to obtain a refinement of the Strichartz/Stein-Tomas estimates for (1.1) for data in a (generalized) hat-Morrey space, which is wider than the above hat-Lebesgue spaces (see Appendix A). Let us first give its definition.

**Definition 1.2** (A Morrey and a hat-Morrey spaces). For $j, k \in \mathbb{Z}$, let $\tau_j^k = [k2^{-j}, (k+1)2^{-j})$ be a dyadic interval.

(i) For $1 \leq \gamma \leq \beta \leq \infty$ and $\beta < \delta \leq \infty$, we define a Morrey norm by

$$\| f \|_{M_{\gamma, \delta}^\beta} = \left\| \| f \|_{L^\gamma(\tau_j^k)} \right\|_{\ell^\beta_{j,k}},$$

where, the case $\beta = \gamma$ and $\gamma < \infty$ is excluded.

(ii) For $1 \leq \beta \leq \gamma \leq \infty$ and $\beta' < \delta \leq \infty$, we define a hat-Morrey norm by

$$\| f \|_{\hat{M}_{\gamma', \delta}^\beta} := \| \hat{f} \|_{\hat{M}_{\gamma', \delta}^{\beta'}} = \left\| \| \hat{f} \|_{L^{\gamma'}(\tau_j^k)} \right\|_{\ell^\beta_{j,k}}.$$

Banach spaces $M_{\gamma, \delta}^\beta$ and $\hat{M}_{\gamma', \delta}^\beta$ are defined as sets of tempered distributions of which above norms are finite, respectively.

One of the main motivation of this kind of improvement of the Strichartz estimates lies in its applications to nonlinear theory. Especially, we are interested in construction of a special non-scattering solutions, which are minimal in a suitable sense, to the mass-subcritical generalized Korteweg-de Vries (gKdV) equation:

$$\begin{cases}
\partial_t u + \partial_x^3 u = \mu \partial_x (|u|^{2\alpha} u), & t, x \in \mathbb{R}, \\
u(t_0, x) = u_0(x), & x \in \mathbb{R},
\end{cases}$$

where $t_0 \in \mathbb{R}$, $u : \mathbb{R} \times \mathbb{R} \to \mathbb{R}$ is an unknown function, $u_0 : \mathbb{R} \to \mathbb{R}$ is a given data, and $0 < \alpha < 2$, $\mu \in \mathbb{R}\setminus\{0\}$ are constants. We construct the minimal solution for (gKdV) by using the concentration compactness argument by Kenig-Merle [14]. As explained in [24] Section 1], a good well-posedness theory and a decoupling (in)equality play a central role in the concentration compactness argument. However, when $\alpha < 2$, it seems difficult to derive those properties in the Sobolev or hat-Lebesgue spaces by several reasons. In [24], it turns out that a use of the generalized hat-Morrey space enables us to establish well-posedness theory good enough and to obtain the concentration compactness lemma equipped with a decoupling inequality. Our estimate

---

1Although the decoupling (in)equality is not obtained in $\hat{L}^\alpha$ with $\alpha \neq 2$, it can be established in a wider space $\hat{M}_{2, \delta}^\beta$ thanks to the local $L^2$ structure of $\hat{M}_{2, \delta}^\beta$. 
in Theorem 1.3 removes several technical restrictions made in [24]. See Subsection 1.2 below for more details.

As far as the authors know, the refinement of the Stein-Tomas estimate in this direction first appeared in [3] in a context of Schrödinger equation. Besides its own interests, the refined estimate has been studied rather because of its application. In [5], Bourgain use the refined estimate to show a concentration phenomenon of blow-up solutions for the two dimensional mass-critical nonlinear Schrödinger equation. After Bourgain, the refinement of Strichartz estimates are being used, for instance, in the estimate for the maximal function associated to the Schrödinger equation (Moyua, Vargas and Vega [26, 27]), or in the linear profile decomposition in $L^2$-framework for the Schrödinger equation (Merle and Vega [29], Carles and Keraani [6], and Béjour and Vargas [1]). As for the Airy equation (1.1), Kenig, Ponce and Vega [17] showed the the refined estimate and applied it to a study of a concentration of blow-up solution for the mass-critical generalized KdV equation. By using the estimate, Shao [28] proved the linear profile decomposition for Airy equation in $L^2$-framework.

In all above studies, the refinements were restricted to the case $\alpha = 2$ and the diagonal case $p = q$. In [24] the authors proved the refined Stein-Tomas estimate for (1.1) in the case $\alpha \neq 2$ and used it for proving existence of a minimal non-scattering solution for the mass-subcritical generalized KdV equation in the $\hat{M}^p_\gamma$-framework. A similar refinement in the Schrödinger case was done by the first author [22], including its application to existence of a minimal non-scattering solution for the mass-subcritical nonlinear Schrödinger equation in $\hat{M}^p_\gamma$-framework. However, the refinement is still restricted to the diagonal case $p = q$.

Main purpose of this paper is to extend the refinement to the non-diagonal case $p \neq q$, that is, we show refined Strichartz estimates in our terminology. Our main theorems are as follows. We first give the estimate of the Airy equation in the space $L^p_x(R; L^q_t(R))$.

**Theorem 1.3** (Refined Strichartz’ estimate I). Let $\sigma \in (0, 1/4)$. Let $(p, q)$ satisfy

$$0 \leq \frac{1}{p} \leq \frac{1}{4} - \sigma, \quad \frac{1}{q} \leq \frac{1}{2} - \frac{1}{p} - \sigma.$$ 

Define $\alpha$ and $s$ by

$$\frac{2}{p} + \frac{1}{q} = \frac{1}{\alpha}, \quad s = - \frac{1}{p} + \frac{2}{q}.$$ 

Further, we define $\beta$, $\gamma$, and $\delta$ by

$$\frac{1}{\beta} = \frac{1}{\alpha} + \sigma, \quad \frac{1}{\gamma} = \begin{cases} \frac{1}{\beta} - \frac{1}{p} & \text{if } \frac{1}{q} \geq \frac{1}{p} + \sigma, \\ \frac{1}{\beta} - \frac{1}{q} + \sigma & \text{if } \frac{1}{q} < \frac{1}{p} + \sigma, \end{cases} \quad \frac{1}{\delta} = \frac{1}{2} - \frac{1}{\max(p, q)}.$$ 

Then, there exists a positive constant $C$ depending on $p, q, \sigma$ such that the inequality

$$\left\| \partial_x^s e^{-\lambda \partial_x^2} f \right\|_{L^p_x(R; L^q_t(R))} \leq C \left\| \partial_x^s f \right\|_{\hat{M}^p_{\gamma, \beta}}$$

in (1.3)
holds for any $f \in |\partial_x|^{-\sigma} \hat{M}_{\gamma,\delta}^\beta$.

Remark 1.4. For the diagonal case $p = q$, the inequality $\text{(1.3)}$ holds for $\sigma = 0$, see [24, Theorem B.1].

Next we give the estimate for the Airy equation in the space $L^p_t(R; L^q_x(R))$.

Theorem 1.5 (Refined Strichartz’ estimate II). Let $\sigma \in (0, 1/4)$. Let $(p, q)$ satisfy

$$0 \leq \frac{1}{p} \leq \frac{1}{4}, \quad 1 - \frac{1}{p} = \frac{1}{2} - \frac{1}{q} - \sigma.$$

Define $\alpha$ by

$$\frac{2}{p} + \frac{1}{q} = \frac{1}{\alpha}.$$

Further, we define $\gamma$ and $\delta$ by

$$\frac{1}{\gamma} = \begin{cases} \frac{1}{\alpha} - \frac{1}{p} + \sigma & \text{if } \frac{1}{q} \geq \frac{1}{p} - \sigma, \\ \frac{1}{\alpha} - \frac{1}{q} & \text{if } \frac{1}{q} < \frac{1}{p} - \sigma, \end{cases} \quad \frac{1}{\delta} = \frac{1}{\alpha} - \frac{1}{2 \max(p, q)}.$$

Then, there exists a positive constant $C$ depending on $p, q$ such that the inequality

$$\left\| |\partial_x|^\frac{1}{p} e^{-\omega^2 t} f \right\|_{L^p_t(R; L^q_x(R))} \leq C \|f\| \hat{M}_{\gamma,\delta}^\alpha (1.4)$$

holds for any $f \in \hat{M}_{\gamma,\delta}^\alpha$.

We briefly outline the proofs for Theorems $\text{1.3}$, the proof of Theorem $\text{1.5}$ is similar. The diagonal case $p = q$ can be handled by the bilinear technique as in [28, 24]. However, this approach does not work well in the non-diagonal case. Furthermore, due to lack of an interpolation between the Morrey space and the Lebesgue space, the desired estimate does not follow by a simple interpolation. To overcome those difficulties, we take another approach which is based on [1, 21, 24, 34]. As in the diagonal case, we first rewrite the square of the left hand side of $\text{(1.3)}$ into a bi-linear oscillatory integral. We then split the domain of spacetime integral into infinitely many rectangles by a Whitney type decomposition. By the decomposition, the bilinear form is rewritten as the infinite sum of the bilinear forms of which Fourier supports are compact and do not intersect each other. To justify the above decomposition in $L^p_t L^q_x$ or $L^p_t L^2_x$, we have to add a small margin to each rectangles in the Whitney decomposition in purpose of smooth cutoff. Obviously, this margin produces many doublings which disturb orthogonality of the forms. However, if the margin is putted so nicely that the resulting doubling is acceptable then we obtain the desired estimate. The property is summarized as an almost orthogonal property of the Fourier supports of the forms.

In the Schrödinger case, we can put such margin so that the almost orthogonal property is valid (see [1]). However, in the Airy case, the cubic dispersion makes the situation much worse and it seems there is no way to put margin necessary for smooth cutoff. An idea here is to put the margin only in time direction. Although this requires an unpleasant restriction.
in

For any compact subinterval $(1.5)$ the well-posedness of $(gKdV)$ in the scale critical $\hat{\sigma}$. Application 1 – well-posedness for generalized KdV equation.

Local and global well-posedness of the Cauchy problem $(gKdV)$ in a scale critical or subcritical Sobolev space $H^s(\mathbb{R})$, $s \geq s_\alpha$ has been studied by many authors, where $s_\alpha$ is a scale critical exponent, i.e., $s_\alpha := 1/2 - 1/\alpha$. A fundamental work on local well-posedness is due to Kenig, Ponce and Vega [10]. They proved that $(gKdV)$ is locally well-posed in $H^s(\mathbb{R})$ with $s > 3/4 (\alpha = 1/2)$, $s \geq 1/4 (\alpha = 1)$, $s \geq 1/12 (\alpha = 3/2)$ and $s \geq s_\alpha (\alpha \geq 2)$. Furthermore, in [10] Kenig, Ponce and Vega proved the small data global well-posedness and scattering of $(gKdV)$ in the scale critical space $\hat{H}^{s_\alpha}$ for $\alpha \geq 2$. Tao [30] proved global well-posedness for small data for $(gKdV)$ with the quartic nonlinearity $\mu \partial_x(u^4)$ in $\hat{H}^{s_\alpha/2}$, see also Koch and Marzuola [19] for the simplified proof of [30] and made an extension. Recently, the authors [23] obtained global well-posedness for small data for $(gKdV)$ in the scale critical space $\hat{L}^\alpha$ with $8/5 < \alpha < 10/3$.

We consider global well-posedness for small data for $(gKdV)$ in a scale critical hat-Morrey space $|\partial_x|^{-\sigma} M_{\gamma,\delta}^\beta$ space. It is known that the nonlinear Schrödinger equation is globally well-posed for small data in the scale critical $M_{\gamma,\delta}^\beta$ space, see [22, 24]. In the Schrödinger case, we only need the diagonal refined estimate to obtain well-posedness. On the other hand, as for $(gKdV)$, due to the presence of derivatives in the nonlinearity, we also need the non-diagonal refined Strichartz estimate for $(1.1)$ to yield a similar well-posedness result. In this paper, by using the refined non-diagonal estimate in Theorem 1.3, we shall prove the global well-posedness for small data for $(gKdV)$ in the scale critical $\hat{M}_{\gamma,\delta}^\beta$ space.

Assumption 1.6. Let $5/3 < \alpha \leq 20/9$ and $0 < \sigma \leq \min(3/5 - 1/\alpha, 1/4 - 2/(5\alpha))$. Define $\beta$ by $1/\beta = 1/\alpha + \sigma$. Let $\gamma$ and $\delta$ satisfy

$$\frac{4}{5\alpha} + 2\sigma \leq \frac{1}{\gamma} < \frac{1}{\beta}, \quad \frac{1}{2} - \frac{1}{5\alpha} \leq \frac{1}{\delta} < \frac{1}{\beta'}.$$

Theorem 1.7 (Local well-posedness in $|\partial_x|^{-\sigma} M_{\gamma,\delta}^\beta$). Suppose $\alpha, \sigma, \beta, \gamma,$ and $\delta$ satisfy Assumption 1.6. Then, the initial value problem $(gKdV)$ is locally well-posed in $|\partial_x|^{-\sigma} M_{\gamma,\delta}^\beta$. More precisely, for any $|\partial_x|^{-\sigma} u_0 \in M_{\gamma,\delta}^\beta(\mathbb{R})$, there exist an interval $I = I(u_0)$ and a unique solution to $(gKdV)$ satisfying

$$(1.5) u \in C(I; |\partial_x|^{-\sigma} M_{\gamma,\delta}^\beta(\mathbb{R})) \cap L^4_{dx}(\mathbb{R}; L^5_{t}(I)) \cap |\partial_x|^{-\sigma} L^{3\beta}_{t,x}(I \times \mathbb{R}).$$

For any compact subinterval $I' \subset I$, there exists a neighborhood $V$ of $u_0$ in $|\partial_x|^{-\sigma} M_{\gamma,\delta}^\beta(\mathbb{R})$ such that the map $u_0 \mapsto u$ from $V$ into the class defined by $(1.5)$ with $I'$ instead of $I$ is Lipschitz continuous. The solution satisfies $u(t) - e^{-(t-t_0)|\partial_x|^\beta} u(t_0) \in C(I; \hat{L}^\alpha \cap |\partial_x|^{-\sigma} \hat{L}^\beta)$ for any $t_0 \in I$. 

$\sigma > 0$ in Theorems 1.3 and 1.5 we recover the almost orthogonal property. See Proposition 2.1 and Remark 2.2 for the detail.

Next we give several applications of our refinement estimates.

1.1. Application 1 – well-posedness for generalized KdV equation. As the first application of the refinement of Strichartz’ estimates, we show the well-posedness of $(gKdV)$ in the scale critical $\hat{M}_{\gamma,\delta}^\beta$ space.
Throughout this paper, we call a function \( u \) which satisfies (1.5) and solves the corresponding integral equation as a \(|\partial_x|^{-\sigma}\hat{M}_{\gamma,\delta}\)-solution to \((gKdV)\) on an interval \( I \).

**Theorem 1.8** (Small data scattering in \(|\partial_x|^{-\sigma}\hat{M}_{\gamma,\delta}\)). Suppose \( \alpha, \sigma, \beta, \gamma, \) and \( \delta \) satisfy Assumption 1.6. Then, there exists \( \varepsilon_0 > 0 \) such that if \(|\partial_x|^{\sigma}u_0 \in \hat{M}_{\gamma,\delta}^\beta(\mathbb{R})\) satisfies \( \||\partial_x|^{\sigma}u_0\|_{\hat{M}_{\gamma,\delta}^\beta} \leq \varepsilon_0 \), then the solution \( u(t) \) to \((gKdV)\) given in Theorem 1.7 is global in time and scatters for both time directions. Moreover,

\[
\|\partial_x|^{\sigma}u\|_{L^\infty(\mathbb{R};\hat{M}_{\gamma,\delta}^\beta)} + \|u\|_{L^5(\mathbb{R};L^{5/2}(\mathbb{R}))} \leq 2\||\partial_x|^{\sigma}u_0\|_{\hat{M}_{\gamma,\delta}^\beta}.
\]

**Remark 1.9.** For scale subcritical spaces, there are many results on the small data scattering for the generalized KdV equation \((gKdV)\) for \( \alpha \geq 1 \), see \cite{8, 9, 12} for instance.

1.2. **Application 2 – existence of minimal non-scattering solution.** We next apply the refined Strichartz estimate to construct a minimal non-scattering solution to mass-subcritical generalized KdV equation \((gKdV)\).

As for \((gKdV)\), the mass-critical case \( \alpha = 2 \) is most extensively studied in this direction. Killip-Kwon-Shao-Visan \cite{18} constructed a minimal blow-up solution to the mass critical KdV equation with the focusing nonlinearity in the framework of \(L^2\). Dodson \cite{10} proved the global well-posedness and scattering in \(L^2\) for the mass critical KdV equation with the defocusing nonlinearity.

The authors \cite{24} showed a existence of a minimal non-scattering solution of \((gKdV)\) with the mass-subcritical case. We constructed the critical element by establishing the concentration compactness in the framework of \(\hat{M}_{2,\delta}^\alpha\) space. On the other hand, well-posedness result was not proved in \(\hat{M}_{2,\delta}^\alpha\) but in \(\tilde{L}^\alpha\) due to lack of non-diagonal refined Strichartz estimate. This disagreement caused some technical restrictions in the previous result \cite{24}.

In this paper, by using the non-diagonal refined Strichartz estimate (Theorem 1.3), we resolve the disagreement and show existence of critical element under a reasonable assumption.

Before we state our main theorems in this subsection, we introduce several notation. In the rest of this section, a solution always implies a \(|\partial_x|^{-\sigma}\hat{M}_{2,\delta}^\beta\)-solution unless otherwise stated. We introduce a deformations associated with the function space \(|\partial_x|^{-\sigma}\hat{M}_{2,\delta}^\beta\):

- Translation in Physical side: \((T(y)f)(x) := f(x - y)\).
- Airy flow: \((A(t)f)(x) = (e^{-t\partial_x^3}f)(x)\).
- Dilation (scaling): \((D(h)f)(x) = h^\alpha f(hx)\).

Note that \(|\partial_x|^{-\sigma}\hat{M}_{2,\delta}^\beta\)-norm is invariant under the above group actions.

For a solution \( u \) on \( I \), take \( t_0 \in I \) and set

\[
T_{\max} := \sup\{ T > t_0 \mid u(t) \text{ can be extended to a solution on } [t_0, T) \},
\]

\[
T_{\min} := \sup\{ T > -t_0 \mid u(t) \text{ can be extended to a solution on } (-T, t_0) \},
\]

\[
I_{\max}(u) := (-T_{\min}, T_{\max}).
\]
Definition 1.10 (Scattering). We say a solution $u(t)$ scatters forward in time (resp. backward in time) if $T_{\min} = \infty$ (resp. $T_{\max} = \infty$) and if $|\partial_x|^\sigma e^{it\partial_x^3} u(t)$ converges in $\dot{M}^\beta_{2,\delta}$ as $t \to \infty$ (resp. $t \to -\infty$).

We define

$$E_1 := \inf \left\{ \inf_{t \in I_{\max}} \| |\partial_x|^\sigma u(t) \|_{\dot{M}^\beta_{2,\delta}} : u(t) \text{ is a solution to } \text{[gKdV]} \text{ that does not scatter forward in time.} \right\}.$$ 

Theorem $\text{[LS]}$ is represented as $E_1 > 0$. Remark that it holds that

$$E_1 = \inf \left\{ \| |\partial_x|^\sigma u(0) \|_{\dot{M}^\beta_{2,\delta}} : u(t) \text{ is a solution to } \text{[gKdV]} \text{ that does not scatter forward in time, } 0 \in I_{\max}(u). \right\}$$

by the time translation symmetry. Further, one sees that $E_1$ is the supremum of the number $\epsilon_0$ for which Theorem $\text{[LS]}$ is true.

We also introduce another infimum value.

$$E_2 := \inf \left\{ \lim_{t \to T_{\max}} \| |\partial_x|^\sigma u(t) \|_{\dot{M}^\beta_{2,\delta}} : u(t) \text{ is a solution to } \text{[gKdV]} \text{ that does not scatter forward in time.} \right\}.$$ 

By definition, $E_1 \leq E_2 \leq \| |\partial_x|^\sigma Q \|_{\dot{M}^\beta_{2,\delta}}$. For another characterization of this quantity, see Remark $\text{[Li 15]}$. The goal is to determine the explicit value of $E_j$ ($j = 1, 2$). Here, we will show that existence of minimizers to both $E_1$ and $E_2$, which would be an important step.

In what follows, we consider the focusing case $\mu = -1$ only. However, the focusing assumption is used only for assuring $E_j$ are finite. Our analysis work also in the defocusing case $\mu = +1$ if we assume $E_j$ are finite.

Assumption 1.11. We suppose Assumption $\text{[L6]}$ with $\gamma = 2$ and exclude the endpoint cases, i.e., Let $5/3 < \alpha < 12/5$ and $\max(0, 1/2 - 1/\alpha) < \sigma < \min(3/5 - 1/\alpha, 1/4 - 2/(5\alpha))$. Define $\beta \in (5/3, 2)$ by $1/\beta = 1/\alpha + \sigma$ and let $1/\delta \in (1/2 - 1/(5\alpha), 1/\beta').$

Theorem 1.12 (Analysis of $E_1$). Suppose that Assumption $\text{[L11]}$ is satisfied. Then, $0 < E_1 \leq c_\alpha \| |\partial_x|^\sigma Q \|_{\dot{M}^\beta_{2,\delta}}$. Furthermore, there exists a minimizer $u_1(t)$ to $E_1$ in the following sense: $u_1(t)$ is a solution to $\text{[gKdV]}$ with maximal interval $I_{\max}(u_1) \ni 0$ and

(i) $u_1(t)$ does not scatter forward in time;

(ii) $u_1(t)$ attains $E_1$ in such a sense that either one of the following two properties holds;

(a) $\| |\partial_x|^\sigma u_1(0) \|_{\dot{M}^\beta_{2,\delta}} = E_1$;

(b) $u_1(t)$ scatters backward in time and $u_{1, -} := \lim_{t \to -\infty} e^{it\partial_x^3} u_1(t)$ satisfies $\| |\partial_x|^\sigma u_{1, -} \|_{\dot{M}^\beta_{2,\delta}} = E_1$.

Remark 1.13. Let us mention the difference between the previous results in $\text{[IS 24]}$. In these papers, a priori knowledge of the relation between value of $E_1$ and the same value for a corresponding nonlinear Schrödinger equation is assumed. In our theorem, we do not need this kind of assumption. The assumption is used to exclude the case where $E_1$ is attained by a sequence of initial data of the form $f(x) \cos(\xi_n x)$ with $\xi_n \to \infty$ as $n \to \infty$. The case may happen because the state spaces used in $\text{[IS 24]}$ are $\tilde{L}^\alpha$, in which the
operation \(e^{ix\xi}\) is unitary. In our case, the state space \(|\partial_x|^{-\sigma}\hat{M}^{2,\beta}_{2,\delta}\) contains derivative and so the above case does not take place. Thus, we do not need the assumption.

**Theorem 1.14** (Analysis of \(E_2\)). Suppose that Assumption 1.11 is satisfied. Then, \(E_1 \leq E_2 \leq \||\partial_x|^\sigma Q\|\hat{M}^{\beta}_{2,\delta}\). Furthermore, there exists a minimizer \(u_2(t)\) to \(E_2\) in the following sense: \(u_2(t)\) is a solution to \((\text{gKdV})\) with maximal interval \(I_{\text{max}}(u_2) \ni 0\) and

(i) \(u_2(t)\) does not scatter forward and backward in time;

(ii) Three quantities

\[
\sup_{t \in \mathbb{R}} |||\partial_x|^{\sigma} u_2(t)|||\hat{M}^{\beta}_{2,\delta}, \quad \lim_{t \to T_{\text{max}}} |||\partial_x|^{\sigma} u_2(t)|||\hat{M}^{\beta}_{2,\delta}, \quad \lim_{t \to t_{\text{min}}} |||\partial_x|^{\sigma} u_2(t)|||\hat{M}^{\beta}_{2,\delta}
\]

are equal to \(E_2\).

(iii) \(u_2(t)\) is precompact modulo symmetries, i.e., there exist a scale function \(N(t) : I_{\text{max}} \to \mathbb{R}_+\) and a space center \(y(t) : I_{\text{max}} \to \mathbb{R}\) such that the set \((D(N(t))T(y(t)))^{-1}u_2(t) \mid t \in I_{\text{max}}\) is precompact.

**Remark 1.15.** We give another characterization of \(E_2\). For \(E \geq 0\), we define

\[
\mathcal{L}(E) := \sup \{ |||u|||_{L^\infty_t (R; L^p_x (I))} \mid u(t) \in C(I; |\partial_x|^{-\sigma}\hat{M}^{\beta}_{2,\delta}) \text{ is a solution to } (\text{gKdV}) \text{ on a compact interval } I, \text{ such that } \max_{t \in I} |||\partial_x|^{\sigma} u(t)|||\hat{M}^{\beta}_{2,\delta} \leq E \}.
\]

Remark that \(\mathcal{L} : [0, \infty) \to [0, \infty]\) is non-decreasing. Then, it holds that \(E_2 = \sup \{E \mid \mathcal{L}(E) < \infty\} = \inf \{E \mid \mathcal{L}(E) = \infty\}\).

The following notation will be used throughout this paper: We use the notation \(A \sim B\) to represent \(C_1 A \leq B \leq C_2 A\) for some constants \(C_1\) and \(C_2\). We also use the notation \(A \lesssim B\) to denote \(A \leq CB\) for some constant \(C\). The operator \(|\partial_x|^s = (-\partial_x^2)^{s/2}\) denotes the Riesz potential of order \(-s\). For \(1 \leq p, q \leq \infty\) and \(I \subset \mathbb{R}\), let us define a space-time norm

\[
|||f|||_{L^p_t L^q_x(I)} = |||f(t, \cdot)|||_{L^p_t (\mathbb{R})} |||f(\cdot, \cdot)|||_{L^q_x (\mathbb{R})},
\]

\[
|||f|||_{L^p_t \dot{L}^q_x(I)} = |||f(\cdot, \cdot)|||_{L^p_t \dot{L}^q_x (\mathbb{R})}.
\]

The rest of the article is organized as follows. In Section 2, we prove Theorems 1.3 and 1.5. In Section 3, we shall show the well-posedness and the small data scattering for \((\text{gKdV})\) (Theorems 1.7 and 1.8) by using the refined Strichartz estimate obtained by Theorem 1.3 and the contraction mapping principle. Finally in Section 4, we construct a minimal non-scattering solution to \((\text{gKdV})\) (Theorems 1.12 and 1.14) by using the concentration compactness. In Appendix, we summarize the embedding properties of the generalized Morrey space.

2. **Proof of Strichartz estimates in the hat-Morrey space**

In this section we derive the refinement version of the Strichartz estimates for solution to \((\text{gKdV})\) (Theorems 1.3 and 1.5) by using the argument used in [1] 21, 24 54. We show the inequality (1.3) only, the proof of (1.3) being similar.
2.1. Whitney decomposition. To show the inequality (1.3), we first reduce the linear form into a bilinear form:

$$\|\partial_x|^s e^{-t\partial^2_x} f\|_{L^2_t L^2_x}^2 = \|\partial_x|^s e^{-t\partial^2_x} f\|_{L^2_t L^2_x}^2.$$  

If $f$ is a real valued function, then

$$(e^{-t\partial^2_x} f)(x) = \sqrt{\frac{2}{\pi}} \text{Re} \left[ \int_0^\infty e^{ix\xi + it\xi^3} \hat{f}(\xi) d\xi \right].$$

Hence

$$\|\partial_x|^s e^{-t\partial^2_x} f\|_{L^2_t L^2_x}^2 = \frac{1}{\pi} \text{Re} \int_0^\infty \int_0^\infty e^{ix(\xi + \eta) + it(\xi^3 + \eta^3)} |\xi\eta|^s \hat{f}(\xi) \hat{f}(\eta) d\xi d\eta \quad + \frac{1}{\pi} \text{Re} \int_0^\infty \int_0^\infty e^{ix(\xi - \eta) + it(\xi^3 - \eta^3)} |\xi\eta|^s \hat{f}(\xi) \hat{\bar{f}}(\eta) d\xi d\eta.$$

We now introduce a Whitney decomposition. Let $D_+ = \{(k2^{-j}, (k+1)2^{-j})| j \in \mathbb{Z}, 0 \leq k \in \mathbb{Z}\}$. For $\tau_k^j, \tau_{\ell}^j \in D_+$, we define a binary relation

$$\tau_k^j \sim \tau_{\ell}^j \iff \begin{cases} \ell - k = -2, 2, 3 & \text{if } k \text{ is even}, \\ \ell - k = -3, -2, 2 & \text{if } k \text{ is odd}. \end{cases}$$

Then, we have $(\mathbb{R}_+ \times \mathbb{R}_+) \setminus \{(\xi, \eta)| \xi \geq 0\} = \bigcup\{\tau_k^j \times \tau_{\ell}^j | \tau_k^j \in D_+, \tau_{\ell}^j \sim \tau_k^j \}$. The Whitney decomposition gives us

$$\|\partial_x|^s e^{-t\partial^2_x} f\|_{L^2_t L^2_x}^2 = \frac{1}{\pi} \sum_{\tau_k^j \in D_+} \sum_{\tau_{\ell}^j: \tau_{\ell}^j \sim \tau_k^j} \text{Re} \int_{\tau_k^j} \int_{\tau_{\ell}^j} e^{ix(\xi + \eta) + it(\xi^3 + \eta^3)} |\xi\eta|^s \hat{f}(\xi) \hat{f}(\eta) d\xi d\eta \quad + \frac{1}{\pi} \sum_{\tau_k^j \in D_+} \sum_{\tau_{\ell}^j: \tau_{\ell}^j \sim \tau_k^j} \text{Re} \int_{\tau_k^j} \int_{\tau_{\ell}^j} e^{ix(\xi - \eta) + it(\xi^3 - \eta^3)} |\xi\eta|^s \hat{f}(\xi) \hat{\bar{f}}(\eta) d\xi d\eta \quad = \quad 2 \text{Re} \sum_{\tau_k^j \in D_+} \sum_{\tau_{\ell}^j: \tau_{\ell}^j \sim \tau_k^j} |\partial_x|^s e^{-t\partial^2_x} f_{\tau_k^j} |\partial_x|^s e^{-t\partial^2_x} f_{\tau_{\ell}^j} \quad + 2 \text{Re} \sum_{\tau_k^j \in D_+} \sum_{\tau_{\ell}^j: \tau_{\ell}^j \sim \tau_k^j} |\partial_x|^s e^{-t\partial^2_x} f_{\tau_k^j} |\partial_x|^s e^{-t\partial^2_x} f_{\tau_{\ell}^j} \quad = : \quad I_1 + I_2,$$

where $\hat{f}(\xi) = 1_{I_1} \hat{f}(\xi)$. A simple calculation leads

$$\text{supp } F_{t,x}[|\partial_x|^s e^{-t\partial^2_x} f_{\tau_k^j} |\partial_x|^s e^{-t\partial^2_x} f_{\tau_{\ell}^j}](\tau, \xi) \quad \subset \quad \{(\xi_1^3 + \xi_2^3, \xi_1 + \xi_2)| \xi_1 \in \tau_k^j, \xi_2 \in \tau_{\ell}^j \} \quad \subset \quad A_{j,k,\ell},$$

where $A_{j,k,\ell}$ is given by

$$\forall (\tau, \xi) \quad A_{j,k,\ell} = \left\{ (\tau, \xi) \mid \frac{k + \ell}{2} \leq \xi \leq \frac{k + \ell + 2}{2}, \tau \right\} \quad \text{satisfies } (2.4).$$
with

\[
\begin{cases}
3 \frac{(k - \ell - 1)^2}{4 2^{2j}} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(k - \ell + 1)^2}{4 2^{2j}} \xi \\
3 \frac{(k - \ell + 1)^2}{4 2^{2j}} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(k - \ell - 1)^2}{4 2^{2j}} \xi
\end{cases}
\]
if \( \ell - k = -3, -2, \)

\[
4 \frac{(k + \ell)^2}{4 2^{2j}} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 4 \frac{(k + \ell + 2)^2}{4 2^{2j}} \xi
\]
if \( \ell - k = 2, 3. \)

In a similar way, we see

\[
\text{supp } \mathcal{F}_{t,x}[|\partial_x|^s e^{-t|\partial_x|^r} f](\tau, \xi)
\subset \{(\xi_1^3 + \xi_2^3, \xi_1 + \xi_2) | \xi_1 \in \tau R, \xi_2 \in \tau_{-\ell-1} R, \lambda > 0, \}
\subset B_{j,k,\ell},
\]
where \( B_{j,k,\ell} \) is given by

\[
B_{j,k,\ell} = \{(\tau, \xi) \mid \frac{k - \ell - 1}{2^j} \leq \xi \leq \frac{k - \ell + 1}{2^j}, \tau \text{ satisfies (2.6)}\}.
\]

with

\[
\begin{cases}
3 \frac{(k + \ell)^2}{4 2^{2j}} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(k + \ell + 2)^2}{4 2^{2j}} \xi \quad \text{if } \ell - k = -3, -2, \\
3 \frac{(k + \ell + 2)^2}{4 2^{2j}} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(k + \ell)^2}{4 2^{2j}} \xi \quad \text{if } \ell - k = 2, 3.
\end{cases}
\]

2.2. Key estimates. Let us introduce two preliminary estimates associated with the set \( A_{j,k,\ell} \) and \( B_{j,k,\ell} \) given in the previous section.

For a closed domain \( R \subset \mathbb{R}^2 \) and \( \lambda > 0 \), we define

\[
R_{+\lambda} = \{(\tau + \tau', \xi) | (\tau, \xi) \in R, -\lambda \leq \tau' \leq \lambda\}.
\]

The set \( R_{+\lambda} \) is an enlargement of \( R \) in \( \tau \)-direction. Let \( \varphi \in C_0^\infty(\mathbb{R}) \) be a nonnegative function such that \( \supp \varphi \subset [-1, 1] \) and \( \int_{-1}^1 \varphi(x) dx = 1 \). Define a cut-off function

\[
\psi_{R,\lambda}(\tau, \xi) := \left[ \frac{2}{\lambda^3} \left( \frac{2}{\lambda} \right) * \tau \right] \left[ 1_{R_{+\lambda}}(\tau', \xi) \right](\tau),
\]

where \( \psi_{R,\lambda} \) is a characteristic function supported on \( \Omega \subset \mathbb{R}^2 \). Note that \( \psi_{R,\lambda} \) is a smooth function with respect to \( \tau \) variable. Furthermore, \( \psi_{R,\lambda} \) satisfies \( 0 \leq \psi_{R,\lambda} \leq 1, \psi_{R,\lambda} \equiv 1 \) on \( R \), and \( \supp \psi_{R,\lambda} \subset R_{+\lambda} \). We define a Fourier multiplier \( P_{R,\lambda} \) by

\[
(P_{R,\lambda} f)(t, x) := \mathcal{F}_{\tau,\xi}^{-1}[\psi_{R,\lambda} \mathcal{F}_{t,x} f](t, x)
\]
\[
= \left( \mathcal{F}_{\tau,\xi}^{-1}[\varphi \left( \frac{2}{\lambda} \right) \tau] \mathcal{F}_{\tau,\xi}^{-1}[1_{R_{+\lambda}}] * f \right)(t, x).
\]

Let \( \Lambda = \{(j, k, \ell) \in \mathbb{Z} \times \mathbb{Z}_{\geq 0} \times \mathbb{Z}_{\geq 0} | \ell - k = -3, -2, 2, 3\} \). For \( (j, k, \ell) \in \Lambda \), we let two families of sets \( \{A_{j,k,\ell}\} \) and \( \{B_{j,k,\ell}\} \) be as in (2.3) and (2.5), respectively. We further introduce

\[
\tilde{A}_{j,k,\ell} = (A_{j,k,\ell})_{+ \frac{1}{100 \times 2^{2j}}}, \quad \tilde{B}_{j,k,\ell} = (B_{j,k,\ell})_{+ \frac{1}{100 \times 2^{2j}}}.
\]
As we explained in Introduction, the following finite doubling properties of the two families \( \{ \tilde{A}_{j,k,\ell} \} \) and \( \{ \tilde{B}_{j,k,\ell} \} \) play an important role in the proof of Theorems 1.3 and 1.5.

**Proposition 2.1** (Almost orthogonality). Let \( X = A \) or \( B \). Then the inequality

\[
\sum_{(j,k,\ell) \in \Lambda} 1_{\tilde{A}_{j,k,\ell}}(\tau, \xi) \leq 12
\]

holds for almost all \((\tau, \xi) \in \mathbb{R}^2\), where \( \Lambda = \{(j, k, \ell) \in \mathbb{Z} \times \mathbb{Z}_{\geq 0} \times \mathbb{Z}_{\geq 0} | \ell - k = -3, -2, 2, 3\} \).

**Proof of Proposition 2.1.** We first note that \( \Lambda = \{(j, k, k + m) \in \mathbb{Z} \times \mathbb{Z}_{\geq 0} \times \mathbb{Z}_{\geq 0} | m = -3, -2, 2, 3\} \). By (2.7), we see

\[
\tilde{A}_{j,k,k+m} \subset \left\{ (\tau, \xi) \middle| \frac{2(k + m)}{2^j} \leq \xi \leq \frac{2(k + 1) + m}{2^j}, \tau \text{ satisfies } (2.10) \right\}
\]

where

\[
(2.10) \begin{cases} 
\frac{3}{16} \frac{(m + 1)^2}{2^j} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(m - 1)^2}{2^j} \xi & \text{if } m = -3, -2, \\
\frac{3}{16} \frac{(m - 1)^2}{2^j} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(m + 1)^2}{2^j} \xi & \text{if } m = 2, 3.
\end{cases}
\]

Hence for any \( j \in \mathbb{Z} \) and \( m = -3, -2, 2, 3 \), we have

\[
\sum_{k \geq \max\{0,-m\}} 1_{\tilde{A}_{j,k,k+m}}(\tau, \xi) \leq 1_{C_{j,m}}(\tau, \xi) \quad a.e.,
\]

where \( C_{j,m} = \{(\tau, \xi) | \xi \geq 0, \tau \text{ satisfies } (2.10) \} \). Therefore

\[
(2.11) \sum_{j \in \mathbb{Z}} \sum_{k \geq \max\{0,-m\}} 1_{\tilde{A}_{j,k,k+m}}(\tau, \xi) \leq 3 \quad a.e.
\]

for each \( m = -3, -2, 2, 3 \). Hence we obtain (2.9) for \( X = A \).

By (2.7), we find

\[
\tilde{B}_{j,k,k+m} \subset \left\{ (\tau, \xi) \middle| \frac{-m - 1}{2^j} \leq \xi \leq \frac{-m + 1}{2^j}, \tau \text{ satisfies } (2.12) \right\}
\]

where

\[
(2.12) \begin{cases} 
\frac{3}{16} \frac{(2k + m)^2}{2^j} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(2k + 1) + m)^2}{2^j} \xi & \text{if } m = -3, -2, \\
\frac{3}{16} \frac{(2k + 1) + m)^2}{2^j} \xi \leq \tau - \frac{1}{4} \xi^3 \leq 3 \frac{(2k + 1)}{2^j} \xi & \text{if } m = 2, 3.
\end{cases}
\]

Therefore, for any \( j \in \mathbb{Z} \) and \( m = -3, -2, 2, 3 \) we have

\[
\sum_{k \geq \max\{0,m\}} 1_{\tilde{B}_{j,k,k+m}}(\tau, \xi) \leq 31_{D_{j,m}}(\tau, \xi) \quad a.e.,
\]

where \( D_{j,m} \) are given by

\[
D_{j,m} = \left\{ (\tau, \xi) \in \mathbb{R}^2 \middle| \frac{-m - 1}{2^j} \leq \xi \leq \frac{-m + 1}{2^j} \right\}.
\]
This implies
\[
\sum_{j \in \mathbb{Z}} \sum_{k \geq \max \{0, -m \}} 1_{\tilde{B}_{j,k,k+m}}(\tau, \xi) \leq 3 \quad \text{a.e.}
\]
for any \( m = -3, -2, 2, 3 \). Then we obtain (2.9) for \( X = B \). This completes the proof of Proposition 2.1. \( \square \)

**Remark 2.2.** For closed domain \( R \subset \mathbb{R}^2 \) and \( \lambda > 0 \), we define
\[
\tilde{R}'_{\lambda} = \{(\tau + \tau', \xi + \xi')| (\tau, \xi) \in R, -\lambda \leq \tau' \leq \lambda, -\lambda \leq \xi' \leq \lambda\},
\]
which is an enlargement both in \( \tau \)- and \( \xi \)-directions. Further, we define \( \tilde{A}'_{j,k,\ell} \) and \( \tilde{B}'_{j,k,\ell} \) by
\[
\tilde{A}'_{j,k,\ell} = (A_{j,k,\ell})'_{\frac{k}{100 \times 2^{2\ell}}}, \quad \tilde{B}'_{j,k,\ell} = (B_{j,k,\ell})'_{\frac{k}{100 \times 2^{2\ell}}}.
\]
If we are able to show the almost orthogonality properties of the two families \( \{\tilde{A}'_{j,k,\ell}\} \) and \( \{\tilde{B}'_{j,k,\ell}\} \), then we will be able to obtain Theorems 1.3 and 1.5 with \( \sigma = 0 \) by means of [31, Lemma 6.1], in essentially the same spirit as in [1].

Next we show the bounds for the Fourier multipliers \( P_{\tilde{A}_{j,k,\ell}} \) and \( P_{\tilde{B}_{j,k,\ell}} \) defined by
\[
P_{\tilde{X}_{j,k,\ell}} := P_{X_{j,k,\ell}}_{\frac{k}{100 \times 2^{2\ell}}} \quad \text{for } X = A, B,
\]
where \( P_{X_{j,k,\ell}}_{\frac{k}{100 \times 2^{2\ell}}} \) is given by (2.7). Since \( P \) is a frequency cutoff which is smooth only in \( \tau \)-direction, we are not free from a small loss in the exponent of \( x \).

**Proposition 2.3** (Bounds for multiplier). Let \( X = A \) or \( B \).

(i) Let \( \sigma > 0 \), \( 1/(1 - \sigma) \leq p \leq \infty \) and \( 1 \leq q \leq \infty \). Let \( p_\sigma \) be given by
\[
\frac{1}{p_\sigma} = \frac{1}{p} + \sigma.
\]
Then, there exists a positive constant \( C \) depending only on \( p, q \) such that for any \((j, k, \ell) \in \Lambda\), the inequality
\[
\|P_{\tilde{X}_{j,k,\ell}} F\|_{L^p_tL^q_x} \leq C 2^{-j\sigma} \|F\|_{L^p_tL^q_x}
\]
holds for any \( F \in L^p_t L^q_x \).

(ii) Let \( \sigma > 0 \), \( 1 \leq p \leq \infty \) and \( 1/(1 - \sigma) \leq q \leq \infty \). Let \( q_\sigma \) be given by
\[
\frac{1}{q_\sigma} = \frac{1}{q} + \sigma.
\]
Then, there exists a positive constant \( C \) depending only on \( p, q \) such that for any \((j, k, \ell) \in \Lambda\), the inequality
\[
\|P_{\tilde{X}_{j,k,\ell}} F\|_{L^p_tL^q_x} \leq C 2^{-j\sigma} \|F\|_{L^p_tL^q_x}
\]
holds for any \( F \in L^p \ L^q_x \).
Proof of Proposition 2.3. We show the inequality (2.15) only since the proof of (2.10) is similar. Consider a set of the form

\[ R := \left\{ (\tau, \xi) \mid a \leq \xi \leq b, c \xi \leq \tau - \frac{1}{4} \xi^3 \leq d \xi \right\} \]

with constants \( a, b, c, d \in \mathbb{R} \). To prove (2.15), it suffices to show the following assertion: For \( a, b, c, d \in \mathbb{R} \) satisfying the relations

\[ (b + a)(d - c) \leq 200 \frac{k}{2^j}, \quad b - a \leq \frac{2}{2^j} \]

and for \( \lambda = k/2^3j \), it holds that

(2.17) \[ \| P_{R, \lambda} F \|_{L^q_x L^q_t} \leq C 2^{-j} \| F \|_{L^q_x L^q_t}. \]

To prove (2.17), we first evaluate the inverse Fourier transform of the characteristic function \( 1_{R \times \frac{1}{2}} \). A direct calculation shows

\[
\mathcal{F}^{-1}[1_{R \times \frac{1}{2}}](t, x) = \int_a^b \left( \int_{\frac{1}{4} \xi^3 + \frac{1}{2} \eta}^{\frac{3}{4} \xi^3 + \frac{1}{2} \eta} e^{i x \xi + i t \tau} \, d\tau \right) \, d\xi
\]

\[
= \int_a^b e^{i(x + f t) \xi + \frac{1}{4} i t \xi^3} \left( \int_{-\frac{1}{2} \xi - \frac{1}{2} t}^{\frac{1}{2} \xi + \frac{1}{2} t} e^{i t \tau} \, d\tau \right) \, d\xi,
\]

where \( f = (d + c)/2 \). We easily see

(2.18) \[ |\mathcal{F}^{-1}[1_{R \times \frac{1}{2}}](t, x)| \leq \int_a^b \left( \int_{-\frac{1}{2} \xi - \frac{1}{2} t}^{\frac{1}{2} \xi + \frac{1}{2} t} \, d\tau \right) \, d\xi
\]

\[ \leq \frac{1}{2} (b^2 - a^2)(d - c) + \lambda(b - a) \leq C \frac{k}{2^j}. \]

On the other hand, we evaluate \( \mathcal{F}^{-1}[1_{R \times \frac{1}{2}}] \) by using the method of stationary phase. We rewrite \( \mathcal{F}^{-1}[1_{R \times \frac{1}{2}}] \) as

\[
\mathcal{F}^{-1}[1_{R \times \frac{1}{2}}](t, x)
\]

\[ = \int_a^b e^{i(x + (\frac{3}{4} e^2 + f)t)\xi + \frac{1}{4} i t (\xi^3 - 3 \xi^2 \xi)} \left( \int_{-\frac{1}{2} \xi - \frac{1}{2} t}^{\frac{1}{2} \xi + \frac{1}{2} t} \frac{d e^{i t \tau}}{d \xi} \, d\tau \right) \, d\xi,
\]

where \( e = (a + b)/2 \). Using the identity

\[ e^{i(x + (\frac{3}{4} e^2 + f)t)\xi} = \frac{\partial e^{i(x + (\frac{3}{4} e^2 + f)t)\xi}}{i(x + (\frac{3}{4} e^2 + f)t)\xi}
\]

and integrating by parts, we obtain

\[
\mathcal{F}^{-1}[1_{R \times \frac{1}{2}}](t, x)
\]

\[ = \frac{1}{i(x + (\frac{3}{4} e^2 + f) t)} \left[ e^{i x \xi + \frac{1}{4} i t \xi^3} \left( \int_{-\frac{1}{2} \xi - \frac{1}{2} t}^{\frac{1}{2} \xi + \frac{1}{2} t} e^{i t \tau} \, d\tau \right) \right]_a^b
\]

\[ - \frac{3 \lambda}{4(x + (\frac{3}{4} e^2 + f) t)} \int_a^b (\xi^2 - e^2) e^{i x \xi + \frac{1}{4} i t \xi^3} \left( \int_{-\frac{1}{2} \xi - \frac{1}{2} t}^{\frac{1}{2} \xi + \frac{1}{2} t} e^{i t \tau} \, d\tau \right) \, d\xi
\]

REFINEMENT OF STRICHARTZ ESTIMATE FOR AIRY EQUATION 13
Young and Minkowski inequalities, we obtain
\[ (2.19) \]
\[
|F^{-1}[1_{R+\frac{3}{2}}](t, x)| \leq C(b + a)(d - c) + (b - a)^3 + \frac{\lambda}{|x + (\frac{3}{4}e^2 + f)t|} \leq \frac{k}{2^{3j}|x + (\frac{3}{4}e^2 + f)t|}.
\]

By (2.18) and (2.19), we have
\[
|F^{-1}[1_{R+\frac{3}{2}}](t, x)| \leq C \frac{k}{2^{3j}} \times \begin{cases} \frac{1}{2^j} & \text{if } |x + (3e^2/4 + f)t| \leq 2^j, \\ \frac{1}{|x + (\frac{3}{4}e^2 + f)t|} & \text{if } |x + (3e^2/4 + f)t| \geq 2^j. \end{cases}
\]

Therefore we find
\[
\|F^{-1}[1_{R+\frac{3}{2}}]\|_{L_t^\infty L_x^q} \leq C \frac{k}{2^{3j}} 2^{-j\alpha},
\]
where \( r_\alpha \) satisfies \( -\alpha = 1/r_\alpha - 1 \). Combining the above inequality with the Young and Minkowski inequalities, we obtain
\[
\|P_{R, \lambda} F\|_{L_t^p L_x^q} \leq C \|F^{-1}[\varphi](\lambda t/2) F^{-1}[1_{R+\frac{3}{2}}]\|_{L_t^p L_x^q} \|F\|_{L_t^p L_x^q} \leq C \|F^{-1}[\varphi](\lambda t/2) F^{-1}[1_{R+\frac{3}{2}}]\|_{L_t^p L_x^q} \|F\|_{L_t^p L_x^q} \leq C 2^{-j\alpha} \|F\|_{L_t^p L_x^q}.
\]

This proves (2.17) and completes the proof. \( \square \)

2.3. Proof of main results. We now prove Theorem 1.3. We may suppose that \( p \neq q \) because the case \( p = q \) is already proved in [24] Theorem B.1. We evaluate \( I_2 \) in (2.2) only since the estimate for \( I_1 \) in (2.2) is similar.

To evaluate \( I_2 \), we first show that
\[
(2.20) \quad \| \sum_{\tau_k^j \in D} \sum_{\tau_k^j \sim \tau_k^j} |\partial_x|^s e^{-it_\ell^j} \partial_{x_k}^s e^{-it_\ell^j} \|_{L_t^p L_x^q} \leq C \left( \sum_{\tau_k^j \in D} \sum_{\tau_k^j \sim \tau_k^j} |\partial_x|^s e^{-it_\ell^j} \partial_{x_k}^s e^{-it_\ell^j} \right)^{\frac{1}{2}},
\]

where the exponent \( p_\sigma \) is given by (2.14). The inequality (2.20) follows from
\[
(2.21) \quad \|TF\|_{L_t^p L_x^q} \leq C \|2^{-j\alpha} F\|_{L_t^p L_x^q}
\]
for function \( F = F(j, k, t, x) \), where \( T \) is defined by
\[
(TF)(t, x) = \sum_{(j, k, t) \in \Lambda} P_{A_{j, k, t}} F(j, k, t, x)
\]
\[
= \sum_{m = -3, -2, 3} \sum_{j > 0} \sum_{k > \max(0, -m)} P_{A_{j, k, k + m}} F(j, k, t, x),
\]
where $\tilde{A}_{j,k,\ell}$ is given by \((2.8)\), $\|a_j\|_{E_j} = (\sum_{j \in \mathbb{Z}} |a_j|^\delta)^{1/\delta}$, and $\|a_k\|_{E_k} = (\sum_{k \in \mathbb{Z}_+} |a_k|^\delta)^{1/\delta}$. Indeed, taking

$$F(j,k,t,x) = \sum_{\tau_1':\tau_1':(j,k)} |\partial_x| e^{-i\tau_1' f_{j,k}} |\partial_x| e^{-i\tau_1 f_{j,k}}$$

in \((2.24)\) and using triangle inequality, we have \((2.23)\).

Let us prove \((2.24)\). The Plancherel identity and the almost orthogonality (Proposition \(2.1\)) imply

$$\|TF\|_{L_t^p L_x^q} \leq C\|F\|_{\dot{E}_{j,k}^p L_t^q L_x^q},$$

On the other hand, the triangle inequality and Proposition \(2.1\) yield

$$\|TF\|_{L_t^p L_x^q} \leq C\|2^{-2j\delta} F\|_{\dot{E}_{j,k}^p L_t^q L_x^q},$$

where

\[
\left(\theta, \frac{1}{P}, \frac{1}{Q}, \frac{1}{P_\sigma}\right) = \begin{cases} \left(\frac{p}{p-4}, 0, \frac{2(p-q)}{q(p-4)}, \frac{2p}{p-4}\right), & \text{if } p > q, \\ \left(\frac{q}{q-4}, \frac{2(q-p)}{p(q-4)}, 0, \frac{2q}{q-4}\right), & \text{if } p < q. \end{cases}
\]

Interpolating \((2.22)\) and \((2.23)\), we obtain \((2.21)\) with

$$\frac{1}{\delta} = 1 - \min\left(\frac{1}{p}, \frac{1}{q}\right) = \frac{1}{2} - \frac{1}{\max(p,q)}.$$ 

To show \((1.3)\), we consider the two cases: $p < q$ or $p > q$.

**Case:** $p < q$. Since $\tau_1' \sim \tau_1^j$ implies $k \neq 0$ or $\ell \neq 0$, we may assume $\ell \neq 0$ in which case $\text{dist}(0, \tau_1^1) \neq 0$.

By an argument used in \cite{24} Proposition B.1, we find

\begin{align*}
(2.24) \quad & \|\partial_x|^{\frac{1}{p} + \sigma} e^{-i\tau_1' f_{j,k}} |\partial_x|^{\frac{1}{p} - \sigma} e^{-i\tau_1 f_{j,k}}\|_{L_t^p L_x^q} \\
& \leq C \text{dist}(0, \tau_1')^{\frac{1}{p} - \sigma} |\tau_1'|^{\frac{1}{p} + \sigma} \|\xi|^{\sigma} \hat{f}_{j,k}\|_{L_t^p L_x^q} \|\xi|^{\sigma} \hat{f}_{j,k}\|_{L_t^p L_x^q},
\end{align*}

where we used the inequality $p_\sigma/2 \geq 2$. Further, 

\begin{align*}
(2.25) \quad & \|\partial_x|^{\frac{1}{p} + \sigma} e^{-i\tau_1' f_{j,k}} |\partial_x|^{\frac{1}{p} - \sigma} e^{-i\tau_1 f_{j,k}}\|_{L_t^p L_x^q} \\
& \leq \|\partial_x|^{\frac{1}{p} - \sigma} e^{-i\tau_1' f_{j,k}}\|_{L_t^p L_x^q} \|\partial_x|^{\frac{1}{p} + \sigma} e^{-i\tau_1 f_{j,k}}\|_{L_t^p L_x^q} \\
& \leq C \text{dist}(0, \tau_1')^{\frac{1}{p} + \sigma} \|\xi|^{\sigma} \hat{f}_{j,k}\|_{L_t^p L_x^q} \|\xi|^{\sigma} \hat{f}_{j,k}\|_{L_t^p L_x^q},
\end{align*}

where we also used the inequality $p_\sigma \geq 4$. Hence, it holds from the Stein interpolation for mixed norm (see \cite{2} Section 7, Theorem 1)), \((2.24)\) and \((2.25)\) that

\begin{align*}
(2.26) \quad & \|\partial_x|^{\frac{1}{p} + \sigma} e^{-i\tau_1' f_{j,k}} |\partial_x|^{\frac{1}{p} - \sigma} e^{-i\tau_1 f_{j,k}}\|_{L_t^p L_x^q} \\
& \leq C |\tau_1'|^{\frac{1}{p} + \sigma} \|\xi|^{\sigma} \hat{f}_{j,k}\|_{L_t^p L_x^q} \|\xi|^{\sigma} \hat{f}_{j,k}\|_{L_t^p L_x^q}.
\end{align*}
Collecting (2.20) and (2.26), we obtain
\[(2.27)\]
\[
\|I_2\|_{L^p_t L^q_x} \leq C\|\partial_x |\sigma f|_2^{\frac{2}{M_{\beta,\gamma}}}. 
\]

**Case:** \( p > q\). As in the previous case, we may assume \( \ell \neq 0 \). By an argument used in \([24]\) Proposition B.1, we find
\[(2.28)\]
\[
\|\partial_x |e^{-t\alpha 2} \sum_{j=1}^{\infty} |\sigma e^{-t\alpha 2} f_j| \|_{L^p_t L^q_x} \leq C \text{dist}(0, \tau_\ell)^{\frac{3}{2} - \frac{2 \rho + \alpha}{\rho q}} \tau_{\ell k}^{\frac{3}{2} - \frac{2 \rho + \alpha}{\rho q}} \|\xi|^{\rho} \hat{f}_{\tau_{\ell k}} \|_{L^p_{\ell k} L^q_{\ell k}} \|\xi|^{\rho} \hat{f}_{\tau_{\ell k}} \|_{L^p_{\ell k} L^q_{\ell k}},
\]
where we used the inequality \( p_{\alpha q} (p_{\alpha q} + q) \geq 2 \). On the other hand, \([23]\) Proposition 2.1 yields
\[(2.29)\]
\[
\|\partial_x |e^{-t\alpha 2} \sum_{j=1}^{\infty} |\sigma e^{-t\alpha 2} f_j| \|_{L^p_t L^q_x} \leq C \text{dist}(0, \tau_\ell)^{\frac{3}{2} - \frac{2 \rho + \alpha}{\rho q}} \tau_{\ell k}^{\frac{3}{2} - \frac{2 \rho + \alpha}{\rho q}} \|\xi|^{\rho} \hat{f}_{\tau_{\ell k}} \|_{L^p_{\ell k} L^q_{\ell k}} \|\xi|^{\rho} \hat{f}_{\tau_{\ell k}} \|_{L^p_{\ell k} L^q_{\ell k}},
\]
Combining Stein interpolation for mixed norm with (2.28) and (2.29), we obtain
\[(2.30)\]
\[
\|\partial_x |e^{-t\alpha 2} \sum_{j=1}^{\infty} |\sigma e^{-t\alpha 2} f_j| \|_{L^p_t L^q_x} \leq C \text{dist}(0, \tau_\ell)^{\frac{3}{2} - \frac{2 \rho + \alpha}{\rho q}} \tau_{\ell k}^{\frac{3}{2} - \frac{2 \rho + \alpha}{\rho q}} \|\xi|^{\rho} \hat{f}_{\tau_{\ell k}} \|_{L^p_{\ell k} L^q_{\ell k}} \|\xi|^{\rho} \hat{f}_{\tau_{\ell k}} \|_{L^p_{\ell k} L^q_{\ell k}}.
\]
Collecting (2.20) and (2.30), we obtain (2.27). For the sub-case \( p_{\alpha} < q\), the similar argument as that in the case \( p < q\) yields (2.27).

In a similar way, we obtain
\[(2.31)\]
\[
\|I_1\|_{L^p_t L^q_x} \leq C\|\partial_x |\sigma f|_2^{\frac{2}{M_{\beta,\gamma}}}. 
\]
Combining (2.27) with (2.31), we obtain (1.3). This completes the proof of Theorem 1.3.

### 3. Application to well-posedness

In this section we prove local and global well-posedness for \((\text{gKdV})\) (Theorems 1.4 and 1.8). To this end, we consider integral form of \((\text{gKdV})\):
\[(3.1)\]
\[
\begin{align*}
\begin{array}{c}
\sigma > 0 \text{ and define } \beta \text{ by } 1/\beta = 1/\alpha + \sigma \text{ as in Theorem 1.7. For an interval } I \subset \mathbb{R}, \text{ we introduce function spaces } L(I), M(I), S(I), \text{ and } D_\sigma(I) \text{ as follows:} \\
L(I) & := \left\{ u \in S'(I \times \mathbb{R}) \left| \|u\|_{L(I)} := \left\| \partial_x |\sigma f| \right\|_{L^p_{\infty} L^q(\mathbb{R}; L^p_{x'}(I))} < \infty \right\} \right., \\
M(I) & := \left\{ u \in S'(I \times \mathbb{R}) \left| \|u\|_{M(I)} := \left\| \partial_x |\sigma f| \right\|_{L^p_{\infty} L^q(\mathbb{R}; L^p_{x'}(I))} < \infty \right\} \right. \\
\end{array}
\end{align*}
\]
\[\begin{array}{c}
(\text{gKdV}) \quad u(t) = u_{0} + \mu \int_{t_0}^{t} e^{-(t-s)\alpha} \partial_x^3 u(s) ds.
\end{array}\]
\[ S(I) := \left\{ u \in S'(I \times \mathbb{R}) \mid \|u\|_{S(I)} := \|u\|_{L^5_x(L^5_t(I))} < \infty \right\}, \]

\[ D_\alpha(I) := \left\{ u \in S'(I \times \mathbb{R}) \mid \|u\|_{D_\alpha(I)} := \|\partial_x^{\alpha} u\|_{L_{x}^{3}\left((\mathbb{R}; L_{t}^{\infty}(I))\right)} < \infty \right\}. \]

For an interval \( I \subset \mathbb{R} \), we say a function \( u \in M(I) \cap S(I) \) is a solution to \((gKdV)\) on \( I \) if \( u \) satisfies (3.1) in the \( M(I) \cap S(I) \) sense. We modify a well-posedness result in [23].

**Lemma 3.1.** Let \( 5/3 < \alpha < 20/9 \). Denote by \( Z(I) \) either \( L(I) \) or \( M(I) \). Let \( t_0 \in \mathbb{R} \) and \( I \) be an interval with \( t_0 \in I \). Then, there exists a universal constant \( \delta > 0 \) such that if a tempered distribution \( u_0 \) and an interval \( I \ni t_0 \) satisfy

\[
\eta_0 = \eta_0(I; u_0, t_0) := \left\| e^{-(t-t_0)\partial_x^3} u_0 \right\|_{S(I)} + \left\| e^{-(t-t_0)\partial_x^3} u_0 \right\|_{Z(I)} \leq \delta,
\]

then there exists a unique solution \( u(t) \) on \( I \) to \((gKdV)\) satisfying

\[
\|u\|_{S(I)} + \|u\|_{Z(I)} \leq 2\eta.
\]

Furthermore, the solution satisfies \( u(t) - e^{-(t-t_0)\partial_x^3} u_0 \in C(I; \hat{L}^\alpha) \).

We omit the proof. Instead, we remark that \( L(I) := X(I; 1/\alpha, \alpha) \) and \( M(I) := X(I; 1/(2\alpha), \alpha) \) in the notation of [23] Definition 1.1], and that the pairs \((s, r) = (1/(2\alpha), \alpha), (1/\alpha, \alpha)\) are acceptable and conjugate-acceptable in the sense of [23] Definitions 1.1 and 3.1] as long as \( 10/7 < \alpha < 8/3 \) if \( Z = L \) and \( 3/2 < \alpha < 7/3 \) if \( Z = M \), which are weaker than our assumption.

As a corollary of Theorem 1.3 and Lemma 3.1 we obtain an existence result.

**Corollary 3.2.** Let \( \alpha, \sigma, \beta, \gamma \) and \( \delta \) satisfy the assumption of Theorem 1.3. Then, for any \( u_0 \in |\partial_x|^{-\sigma} \hat{L}^{\beta, \gamma}_x \) and \( t_0 \in \mathbb{R} \) there exists an interval \( I \subset \mathbb{R} \), \( I \ni t_0 \) such that there exists a unique solution \( u(t) \) on \( I \) to \((gKdV)\). The solution belongs to \( C(I; |\partial_x|^{-\sigma} \hat{L}^{\beta, \gamma}_x + \hat{L}^\alpha) \).

**Proof of Corollary 3.2** One sees from Theorem 1.3 that if \( \alpha > 8/5 \) and \( 0 < \sigma \leq 1/4 - 2/(5\alpha) \) then

\[
(3.2) \quad \|e^{-(t-t_0)\partial_x^3} u_0\|_{L_x(R)} + \|e^{-(t-t_0)\partial_x^3} u_0\|_{S_x(R)} \leq C \|\partial_x^{\alpha} u_0\|_{\hat{L}^{\beta, \gamma}_x} < \infty.
\]

Hence, there exists an open neighborhood \( I \subset \mathbb{R} \) of \( t_0 \) such that \( \eta_0(I; u_0, t_0) \leq \delta \), where \( \delta \) and \( \eta_0 \) are defined in Lemma 3.1. Since \( u(t) - e^{-(t-t_0)\partial_x^3} u_0 \in C(I; \hat{L}^\alpha) \) and \( |\partial_x^\sigma e^{-(t-t_0)\partial_x^3} u_0| \in C(I; \hat{L}^{\beta, \gamma}_x) \), we obtain the result.

**Proof of Theorem 1.7.** To prove the theorem, it suffices to show that \( u(t) - e^{-(t-t_0)\partial_x^3} u_0 \in C(I; |\partial_x|^{-\sigma} \hat{L}^{\beta, \gamma}_x) \). We mimic the argument in [23] Proposition 2.5 that

\[
(3.3) \quad \|u\|_{D_\alpha(I)} \leq C \|\partial_x^{\alpha} u_0\|_{\hat{L}^{\beta, \gamma}_x} + C \|u^{2\alpha} u\|_{N_\alpha(I)},
\]

where

\[
\|f\|_{N_\alpha(I)} := \|\partial_x^{\alpha} f\|_{L^p_x(N_\alpha)} L^q(I^{N_\alpha}),
\]
and \((p(N_\sigma), q(N_\sigma))\) is given by
\[
\frac{2}{p(N_\sigma)} + \frac{1}{q(N_\sigma)} = \frac{1}{\beta} + 2, \quad -\frac{1}{p(N_\sigma)} + \frac{2}{q(N_\sigma)} = \frac{1}{3\beta}.
\]
Note that the pair \((s, r) = (1/(3\beta), \beta)\) is acceptable and conjugate-acceptable in the sense of Definitions 1.1 and 3.1 if \(5/3 \leq \beta < 20/9\). To choose such \(\beta\), we need the restrictions \(5/3 < \alpha \leq 20/9\) and \(0 < \sigma \leq 3/5 - 1/\alpha\). We then apply the Leibniz rule for the fractional order derivatives to obtain
\[
(3.4) \quad \|u\|^{2\alpha} \|u\|_{N_\sigma(I)} \leq C\|u\|_{S(I)}^{2\alpha} \|u\|_{D_\sigma(I)}.
\]
We divide \(I\) into subintervals \(\{I_j\}_{j=1}^J\) so that \(\|u\|_{S(I_j)}\) is small. Then, it follows from \((3.3)\) and \((3.4)\) that \(\|u\|_{D_\sigma(I_j)} < \infty\) for each subinterval, showing \(u \in D_\sigma(I)\). Then, we conclude from the inhomogeneous Strichartz’ estimate Proposition 2.5 that
\[
|\partial_x|^\sigma \int_0^t e^{-(t-s)} \partial_x^2 (|u|^{2\alpha} u) ds. \in C(I; \hat{L}^\beta)
\]
This completes Theorem 3.4, since \(\hat{L}^\beta \hookrightarrow \hat{M}^\beta_{\gamma, \delta}\).

\[\square\]

Remark 3.3. Let us summarize our assumption on the local well-posedness result. For the estimate \((3.3)\), we need \(5/3 < \alpha \leq 20/9\) and \(0 < \sigma \leq 3/5 - 1/\alpha\). Further, the restriction \(\sigma \leq 1/4 - 2/(5\alpha)\) comes from \((3.2)\). The possible ranges of \(\gamma, \delta\) are also ruled by \((3.2)\).

3.1. \textbf{Criterion for blowup and scattering.} We next show standard criterion for finite-time blowup and scattering. These are essentially the same as Theorems 1.8 and 1.9. Let \(I_{\max} = (T_{\min}, T_{\max})\) be the maximal interval.

\textbf{Theorem 3.4} (Criterion of finite-time blowup). Suppose \(\alpha, \sigma, \beta, \gamma, \text{ and } \delta\) satisfy Assumption \(\text{L.}\). Let \(u_0 \in |\partial_x|^{-\sigma} \hat{M}^\beta_{\gamma, \delta}\) and let \(u(t)\) be a corresponding solution given in Theorem \(1.7\) with maximal lifespan \(I_{\max} \ni 0\). If \(T_{\max} < \infty\), then \(\lim_{T \uparrow T_{\max}} \|u\|_{S(0, T)} = \infty\). A similar result holds for backward in time.

\textbf{Theorem 3.5} (Characterization of scattering). Suppose \(\alpha, \sigma, \beta, \gamma, \text{ and } \delta\) satisfy Assumption \(\text{L.}\). Let \(u_0 \in |\partial_x|^{-\sigma} \hat{M}^\beta_{\gamma, \delta}\) and let \(u(t)\) be a corresponding solution given in Theorem \(1.7\) with maximal lifespan \(I_{\max} \ni 0\). The following three statements are equivalent:

- \(u(t)\) scatters forward in time in the sense of Definition \(1.10\);
- \(\|u\|_{L^1(0, T_{\max})} < \infty\);
- \(\|u\|_{S(0, T_{\max})} < \infty\).

Further, if either one of the above (hence all of the above) holds then \(e^{it\beta} u(t)\) converges as \(t \to \infty\) in \(\hat{L}^\alpha \cap |\partial_x|^{-\sigma} \hat{L}^\beta\).

3.2. \textbf{Stability estimate.} By a standard argument, we also obtain a stability estimate. To state it, we introduce a function space with the following norm.
\[
\|f\|_{N(I)} = \|\partial_x^j |\partial_x|^\sigma f\|_{L^p_q(\mathbb{R}; L^r_q(I))},
\]
with
\[
\left( \frac{1}{p(N)}, \frac{1}{q(N)} \right) = \left( \frac{1}{p(M)}, \frac{1}{q(M)} \right) + 2\alpha \left( \frac{1}{p(S)}, \frac{1}{q(S)} \right).
\]

**Theorem 3.6.** Suppose \(\alpha, \sigma, \beta, \gamma, \) and \(\delta\) satisfy Assumption 1.6. Let \(I \subset \mathbb{R}\) be an interval containing \(t_0\). Let \(\tilde{u}\) be an approximate solution to (3.1) on \(I \times \mathbb{R}\) in such a sense that
\[
\tilde{u}(t) = e^{-(t-t_0)d_2} \tilde{u}(t_0) + \int_{t_0}^{t} e^{-(t-s)d_2}(\mu \partial_x (|u|^{2\alpha}u)(s) + e(s))ds
\]
holds in \(L(I) \cap S(I)\) for some function \(e \in N(I)\). Assume that \(\tilde{u}\) satisfies
\[
\|\tilde{u}\|_{S(I)} + \|\tilde{u}\|_{M(I)} \leq M,
\]
for some \(M > 0\). Then there exists \(\varepsilon_1 = \varepsilon_1(M) > 0\) such that if
\[
\|e^{-t\partial_x^2}(u(t_0) - \tilde{u}(t_0))\|_{S(I)} + \|e^{-t\partial_x^2}(u(t_0) - \tilde{u}(t_0))\|_{M(I)} + \|e\|_{N(I)} \leq \varepsilon
\]
and \(0 < \varepsilon < \varepsilon_1\), then there exists a solution \(u\) to (3.1) on \(I \times \mathbb{R}\) satisfies
\[
\|u - \tilde{u}\|_{S(I)} + \|u - \tilde{u}\|_{M(I)} \leq C\varepsilon,
\]
\[
\|u\|_{2\alpha} - |\tilde{u}|_{2\alpha} \|N(I) \leq C\varepsilon,
\]
where the constant \(C\) depends only on \(M\). Further, if \(u(t_0) - \tilde{u}(t_0) \in |\partial_x|^{-\sigma} \hat{M}_{\beta, \delta}\) for some \(\tau \in I\) then, it also holds that
\[
\|\partial_x|^{\sigma}(u - \tilde{u})\|_{L^\infty(I; \hat{M}_{\beta, \delta}^{\sigma})} \leq \|\partial_x|^{\sigma}(u(t_0) - \tilde{u}(t_0))\|_{\hat{M}_{\beta, \delta}^{\sigma}} + C\varepsilon.
\]

**Proof of Theorem 3.6.** Once we obtain Theorem 1.3, the proof follows from the standard continuity argument (See [24, Lemma 3.1, Proposition 3.2]). So we omit the detail. \(\square\)

4. **APPLICATION TO A MINIMIZING PROBLEM**

4.1. **Linear profile decomposition in \(|\partial_x|^{-\sigma} \hat{M}_{\beta, \delta}^{\sigma}\).** In this section, we establish the linear profile decomposition. The linear profile decomposition essentially consists of two parts. The first part is concentration compactness and the second part is the inductive procedure to obtain a decomposition.

Let us begin with the concentration compactness part. The hat-Morrey space \(\hat{M}_{\alpha, \gamma}^{\beta}\) is realized as a dual of a Banach space [22, Theorem 2.17]. Therefore, a bounded set of the hat-Morrey space is compact in the weak-* topology.

**Theorem 4.1.** (Concentration compactness in \(|\partial_x|^{-\sigma} \hat{M}_{\beta, \delta}^{\sigma}\). Suppose that \(\alpha > 8/5\) and \(0 < \sigma < 1/4 - 2/(5\alpha)\). Let \(\beta, \gamma, \delta\) satisfy \(1/\beta = 1/\alpha + \sigma, \)

\[
\frac{4}{5\alpha} + 2\sigma < \frac{1}{\gamma} < \frac{1}{\beta}, \quad \text{and} \quad \frac{1}{2} - \frac{1}{5\alpha} < \frac{1}{\delta} < \frac{1}{\beta}\gamma.
\]

Let \(\{u_n\}_n \subset |\partial_x|^{-\sigma} \hat{M}_{\gamma, \delta}^{\beta}\) a bounded sequence,
\[
\|\partial_x|^{\sigma}u_n\|_{\hat{M}_{\gamma, \delta}^{\beta}} \leq M
\]
for some \(M > 0\). If the sequence further satisfies
\[
\|e^{-t\partial_x^2}u_n\|_{L(\mathbb{R}) \cap S(\mathbb{R})} \geq m
\]
for some \( m > 0 \) then there exist such that
\[
|\partial_x|^{\gamma} (T(y_n)^{-1}A(s_n)^{-1}D(N_n)^{-1}u_n) \to |\partial_x|^{\gamma} \psi
\]
as \( n \to \infty \) weakly-* in \( M_{\eta, \delta}^\beta \) with \( \|\psi\|_{M_{\eta, \delta}^\beta} \geq C(M, m) > 0 \).

**Proof.** In this proof, all spacetime integrals are taken in \( \mathbb{R} \times \mathbb{R} \). Since the endpoint cases are excluded, by means of Theorem 1.3 and by interpolation inequality, we see that the assumption (4.2) implies that \( \|\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u_n\|_{L_{t,x}^{3\alpha}} \geq \tilde{m} \) for some \( \tilde{m} = \tilde{m}(\alpha, \sigma, m) > 0 \). Let \( P_N \) be a standard cut-off operator to \( |\xi| \sim N \in 2^\mathbb{Z} \). We now claim the estimate
\[
\left\| |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u \right\|_{L_{t,x}^{3\alpha}} \leq C \left( \sup_{N \in 2^\mathbb{Z}} \left| P_N |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u \right|_{L_{t,x}^{3\alpha}} \right)^{1 - \frac{\sigma}{\alpha}} \|\partial_x|^\sigma u\|_{M_{\eta, \delta}^\beta}^{\frac{\sigma}{\alpha}},
\]
where \( \zeta := \max(\gamma, \delta) \). By the square function estimate, we have
\[
\left\| |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u \right\|_{L_{t,x}^{3\alpha}} \sim \left\| \sum_{N \in 2^\mathbb{Z}} |P_N |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u|^2 \right\|_{L_{t,x}^{3\alpha}}.
\]
We consider only the case \( 6 < 3\alpha \leq 8 \), the other cases are similar. As \( 3\alpha/8 \leq 1 \),
\[
\text{(R.H.S of (4.3))}^{3\alpha} = \int \prod_{k=1}^4 \left( \sum_{N_k \in 2^\mathbb{Z}} |P_{N_k} |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u|^2 \right)^{3\alpha/8} dt dx
\]
\[
\leq C \sum_{N_1 \leq N_2 \leq N_3 \leq N_4} \int \prod_{k=1}^4 |P_{N_k} |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u|^{3\alpha/4} dt dx.
\]
Let \( \eta > 0 \) be a small number and let \( \alpha_1 = (\frac{1}{3\alpha} - \eta)^{-1} \) and \( \alpha_4 = (\frac{1}{3\alpha} + \eta)^{-1} \).
Remark that \( 2 < \zeta < 3\alpha/2 \). It follows from the H"{o}lder inequality that
\[
\int \prod_{k=1}^4 |P_{N_k} |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u|^{3\alpha/4} dt dx
\]
\[
\leq \left( \sup_{N \in 2^\mathbb{Z}} \left| P_N |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u \right|_{L_{t,x}^{3\alpha}} \right)^{3\alpha - \zeta} \prod_{k=1,4} \left| P_{N_k} |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u \right|_{L_{t,x}^{3\alpha/4}}^{\frac{3\alpha}{4}}.
\]
By Theorem 1.3 we have
\[
|P_{N_k} |\partial_x|^{1/3\alpha} e^{-t\partial_x^3} u|_{L_{t,x}^{3\alpha/4}} \leq C N_k^{\frac{1}{2}} \left( \frac{1}{\alpha_k} - \frac{1}{\alpha_k} \right)^{\frac{1}{4}} \left| P_{N_j} |\partial_x|^{\gamma_k} u \right|_{M_{\eta, \delta}^\beta}
\]
\[
\leq C N_k^{\frac{3}{4}} \left( \frac{1}{\alpha_k} - \frac{1}{\alpha_k} \right)^{\frac{1}{4}} \left| P_{N_j} |\partial_x|^{\gamma_k} u \right|_{M_{\eta, \delta}^\beta}
\]
for \( k = 1, 4 \), where \( \sigma_k, \gamma_k, \delta_k \) are chosen by the relations
\[
\frac{1}{\alpha_k} = \frac{1}{\beta} - \sigma_k = \frac{1}{\alpha} + (\sigma - \sigma_k), \quad \frac{1}{\gamma_k} = \frac{1}{\beta} - \frac{1}{3\alpha_k} \leq \frac{1}{\gamma}, \quad \frac{1}{\delta_k} = \frac{1}{2} - \frac{1}{3\alpha_k} \leq \frac{1}{\delta}.
\]
Remark that $\sigma_1 = \sigma + \eta$ and $\sigma_4 = \sigma - \eta$, and so that the choice is possible if $\eta > 0$ is sufficiently small. Put $a_N := \|P_N|\theta u\|_{M^\beta_{\gamma,\delta}}$ for $N \in 2^\mathbb{Z}$. Combining these inequalities, we reach to the estimate

$$\left\| \partial_x \frac{1}{\alpha} e^{-t\alpha} u \right\|_{L^3_{t,x}} \leq C \left( \sup_{N \in 2^\mathbb{Z}} \| P_N |\theta u\|_{L^3_{t,x}} \right)^{3\alpha - \zeta} \times \sum_{N_1 \in \mathbb{N}_4} a_{N_1}^s a_{N_4}^s \left( \frac{N_1}{N_4} \right)^{\frac{4\alpha}{1 + \log N_4}} \left( 1 + \log \frac{N_4}{N_1} \right)^2.$$

Thus, the claim follows because $\|a_N\|_{L^1} \leq C \|\theta u\|_{M^\beta_{\gamma,\delta}}$ by definition of $\zeta$.

By means of the claim, assumption of the theorem implies that there exists a sequence $\{N_n\} \subset 2^\mathbb{Z}$ such that

$$|N_n|^\frac{1}{3\alpha} \left\| P_{N_n} e^{-t\alpha} u_n \right\|_{L^3_{t,x}} \geq C(M,m).$$

As in (4.4),

$$\left\| P_{N_n} e^{-t\alpha} u_n \right\|_{L^3_{t,x}} \leq \left\| P_{N_n} e^{-t\alpha} u_n \right\|_{L^3_{t,x}}^{1-\theta} \left\| P_{N_n} e^{-t\alpha} u_n \right\|_{L^3_{t,x}}^{\theta} \leq C \left\| P_{N_n} e^{-t\alpha} u_n \right\|_{L^3_{t,x}}^{1-\theta} \left( MN_n^{-\frac{4\alpha}{1 + \log N_4}} \right)^\theta,$$

where $\theta = \alpha_4/\alpha = (1 + \eta)a)^{-1}$. We obtain

$$(N_n)^{-\frac{1}{\alpha}} \left| e^{-t\alpha} P_{N_n} u_n \right|_{L^\infty_{t,x}} \geq C(M,m).$$

Set $v_n(x) := (N_n)^{1/\alpha} u_n(N_n,x)$ to obtain $\|P_1 e^{-t\alpha} v_n\|_{L^\infty_{t,x}} \geq C(M,m)$. Hence, there exists $(s_n, y_n) \in \mathbb{R}^2$ such that

$$\left| P_1 e^{-t\alpha} v_n \right|(y_n) \geq C(M,m).$$

Let $\psi \in \{\partial_x \}^{\alpha} \hat{M}_{\gamma,\delta}^\beta$ be a weak-* limit of $T(y_n) e^{s_n\alpha} v_n$ along a subsequence. Then, by a standard argument, we conclude from (4.5) that $\|\partial_x \}^{\alpha} \hat{M}_{\gamma,\delta}^\beta \geq \beta(M,m)$.}

We next move to the main issue of this section, linear profile decomposition. Let us define a set of deformations as follows

$$G := \{D(N)A(s)T(y) \mid \Gamma = (N, s, y) \in 2^\mathbb{Z} \times \mathbb{R} \times \mathbb{R} \}.$$ We often identify $G \in G$ with a corresponding parameter $\Gamma \in 2^\mathbb{Z} \times \mathbb{R} \times \mathbb{R}$ if there is no fear of confusion. Let us now introduce a notion of orthogonality between two families of deformations.

**Definition 4.2.** We say two families of deformations $\{G_n\} \subset G$ and $\{\tilde{G}_n\} \subset G$ are orthogonal if corresponding parameters $\Gamma_n, \tilde{\Gamma}_n \in 2^\mathbb{Z} \times \mathbb{R} \times \mathbb{R}$ satisfies

$$\lim_{n \to \infty} \left( \left| \log \frac{N_n}{N_n} \right| + |s_n - \left( \frac{N_n}{N_n} \right)^{\frac{3}{2}} s_n| + |y_n - \left( \frac{N_n}{N_n} \right) y_n| \right) = +\infty.$$
Theorem 4.3 (Linear profile decomposition in $|\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$). Suppose that $\alpha$, $\beta$, $\gamma$, and $\delta$ satisfy Assumption 1.11. Let $\{u_n\}_n$ be a bounded sequence in $|\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$. Then, there exist $\psi^j \in |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$, $r_n^j \in |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$, and pairwise orthogonal families of deformations $\{G^j_n\}_n \subset G$ ($j = 1, 2, \ldots$) parametrized by $\{\Gamma^j_n = (h^j_n, s^j_n, \gamma^j_n)\}_n$ such that, extracting a subsequence in $n$,

$$\tag{4.8} u_n = \sum_{j=1}^J G^j_n \psi^j + r_n^j$$

for all $J \geq 1$ and

$$\tag{4.9} \lim_{J \to \infty} \lim_{n \to \infty} \left( \left\| |\partial_x|^{1/2} e^{-t\partial_x^3} r_n^j \right\|_{L^2_{\alpha}(\mathbb{R} \times \mathbb{R})} + \left\| e^{-t\partial_x^3} r_n^j \right\|_{L^2(\mathbb{R} \times \mathbb{R})} \right) = 0$$

Moreover, a decoupling inequality

$$\tag{4.10} \lim_{n \to \infty} \left\| |\partial_x|^\sigma u_n \right\|_{\tilde{M}^{\beta}_{2,\delta}} \geq \sum_{j=1}^J \left\| |\partial_x|^\sigma \psi^j \right\|_{\tilde{M}^{\beta}_{2,\delta}} + \lim_{n \to \infty} \left\| r_n^j \right\|_{\tilde{M}^{\beta}_{2,\delta}}$$

holds for all $J \geq 1$. Furthermore, if $u_n$ is real-valued then so are $\psi^j$ and $r_n^j$.

Proof of Theorem 4.3. For a sequence $\{u_n\}_n \subset |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$, define

$$\mathcal{M}(\{u_n\}) := \left\{ \psi \in |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\sigma} \mid \exists G_n \in G, \exists n_k: \text{subsequence s.t.} \right\}$$

and

$$\eta(\{u_n\}) := \sup_{\phi \in \mathcal{M}(\{u_n\})} \left\| |\partial_x|^\beta \phi \right\|_{\tilde{M}^{\beta}_{2,\delta}}.$$}

Arguing as in [24], one obtains the desired decomposition expect that the smallness (4.9) is replaced by $\tilde{M}^{\beta}_{2,\delta}$. By time translation symmetry, we may suppose that $t_n \equiv 0$. We apply the linear profile decomposition theorem (Theorem 4.3) to the sequence $\{u_n(0)\}_n$. Then, up to subsequence, we obtain a decomposition

$$\tag{4.11} u_n(0) = \sum_{j=1}^J G^j_n \psi^j + r_n^j$$

for $n, J \geq 1$ with the properties (4.9), (4.10), and pairwise orthogonality of $\{G^j_n\}_n \subset G$. By extracting subsequence and changing notations if necessary,
we may assume that for each \( j \) and \( \{x_n^j\}_{n,j} = \{\log N_n^j\}_{n,j}, \{s_n^j\}_{n,j}, \{y_n^j\}_{n,j} \),
either \( x_n^j \equiv 0 \), \( x_n^j \to \infty \) as \( n \to \infty \), or \( x_n^j \to -\infty \) as \( n \to \infty \) holds. Let us define a nonlinear profile \( \Psi^j(t) \) associated with \( (\psi^j, s_n^j) \) as follows: For each \( j \), we let

- if \( s_n^j \equiv 0 \) then \( \Psi^j(t) \) is a solution to \( \text{gKdV} \) with \( \Psi^j(0) = \psi^j \);
- if \( s_n^j \to \infty \) as \( n \to \infty \) then \( \Psi^j(t) \) is a solution to \( \text{gKdV} \) that scatters forward in time to \( e^{-t\partial_x^3} \psi^j \);
- if \( s_n^j \to -\infty \) as \( n \to \infty \) then \( \Psi^j(t) \) is a solution to \( \text{gKdV} \) that scatters backward in time to \( e^{-t\partial_x^3} \psi^j \);

Let
\[
(4.12) \quad V_n^j(t) := D(N_n^j)T(y_n^j)\Psi^j((N_n^j)^3t + s_n^j).
\]
Here, we define an approximate solution
\[
(4.13) \quad \tilde{u}_n^J(t, x) = \sum_{j=1}^J V_n^j(t, x) + e^{-t\partial_x^3}r_n^J.
\]

The main step is to show that there exists \( \Psi^j \) that does not scatter forward in time. Suppose not. Then, all \( \Psi^j \) scatters forward in time and so \( \|\partial_x |s^j| \|_{L^2} < E_1 \) for all \( j \). Then, there exists \( M > 0 \) such that
\[
(4.14) \quad \|V_n^j\|_{L(\mathbb{R}^+) + \|V_n^J\|_{S(\mathbb{R}^+)} \leq M
\]
holds for any \( j, n \geq 1 \).

We shall observe that \( \tilde{u}_n^J \) is an approximately solves \( \text{gKdV} \) and that is close to \( u_n \). To this end, we provide three intermediate results.

**Proposition 4.5** (Asymptotic agreement at the initial time). Let \( \tilde{u}_n^J \) be given by (4.13). Then, it holds for any \( J \geq 1 \) that
\[
\|\tilde{u}_n^J(0) - u_n(0)\|_{L^2} \to 0
\]
as \( n \to \infty \).

**Proposition 4.6** (Uniform bound on the approximate solution). There exists \( M > 0 \) such that
\[
(4.15) \quad \|\tilde{u}_n^J\|_{L(\mathbb{R}^+) + \|\tilde{u}_n^J\|_{S(\mathbb{R}^+)} \leq M
\]
holds for any \( J \geq 1 \) and \( n \geq N(J) \).

**Proposition 4.7** (Approximate solution to the equation). Let \( \tilde{u}_n^J \) be defined by (4.13). Then \( \tilde{u}_n^J \) is an approximate solution to \( \text{gKdV} \) in such a sense that
\[
\lim_{J \to \infty} \lim_{n \to \infty} \sup \|\partial_x^{-1}[(\partial_t + \partial_{xxx})\tilde{u}_n^J - \mu \partial_x(|\tilde{u}_n^J|^{2a} \tilde{u}_n^J)]\|_{N(\mathbb{R}^+)} = 0.
\]

The proof of Proposition 4.5 is obvious by definition of \( V_n^J \). We shall prove Proposition 4.6 later since our improvement is in this proposition. Proposition 4.7 follows from Proposition 4.6 and the mutual asymptotic orthogonality of nonlinear profiles as in [24, Lemma 4.8].
By means of a stability estimate, the above three propositions imply that 
\[ \|u_n\|_{S(\mathbb{R}^+)} < \infty \] for sufficiently large \( n \). This contradicts with the definition of \( \{u_n\}_n \).

Thus, we see that there exists \( j_0 \) such that \( \Psi^{j_0} \) does not scatter. Then, 
\[ |||\partial_x^\sigma \Psi^{j_0}|||_{\tilde{M}_{2,\delta}^3} \geq E_1 \] by definition of \( E_1 \). One also sees from (4.10) that
\[ |||\partial_x^\sigma \Psi^{j_0}|||_{\tilde{M}_{2,\delta}^3} \leq E_1. \] Hence, \[ |||\partial_x^\sigma \Psi^{j_0}|||_{\tilde{M}_{2,\delta}^3} = E_1. \]

Let us show that \( u_c := \Psi^{j_0} \) attains \( E_1 \). The case \( s_n \rightarrow \infty \) as \( n \rightarrow \infty \) is excluded since this implies \( u_c(t) \) scatters forward in time. If \( s_n \equiv 0 \) then \( \Psi^{j_0}(0) = \psi^{j_0} \) and so \[ |||\partial_x^\sigma u_c(0)|||_{\tilde{M}_{2,\delta}^3} = E_1. \] Finally, if \( s_n \rightarrow -\infty \) as \( n \rightarrow \infty \) then \( \lim_{t \to -\infty} e^{t\partial_3^3} \Psi^{j_0}(t) = \psi^{j_0} \). Hence, \( u_{c,-} := \lim_{t \to -\infty} e^{t\partial_3^3} \Psi^{j_0}(t) \) satisfies \[ |||\partial_x^\sigma u_{c,-}|||_{\tilde{M}_{2,\delta}^3} = E_1. \]

To complete the proof of Theorem 1.12 we prove Proposition 4.6. Recall that we have uniform bound (4.14) for each \( V_n^2 \).

**Lemma 4.8.** For any \( \varepsilon > 0 \), there exists \( J_0 = J_0(\varepsilon) \) such that
\[ \left\| \sum_{j=J_0}^{J_0+k} e^{-t\partial_3^3} V_n^j(0) \right\|_{L(\mathbb{R}^+)} + \left\| \sum_{j=J_0}^{J_0+k} e^{-t\partial_3^3} V_n^j(0) \right\|_{S(\mathbb{R}^+)} \leq \varepsilon \]
for any \( k \geq 1 \) and \( n \geq N(k) \).

**Proof.** By definition of \( V_n^2(t) \), it suffices to prove the estimate for \( e^{-t\partial_3^3} G_n^j \psi^j \) instead of \( e^{-t\partial_3^3} V_n^j(0) \). By Theorem 1.3
\[ (4.16) \left\| \sum_{j=J_0}^{J_0+k} e^{-t\partial_3^3} G_n^j \psi^j \right\|_{L(\mathbb{R}^+)} \leq C \left\| \sum_{j=J_0}^{J_0+k} |\partial_x^\sigma G_n^j \psi^j|_{\tilde{M}_{2,\delta}^3} \right\|_{L(\mathbb{R}^+)} \cap S(\mathbb{R}^+). \]

By pairwise orthogonality of \( \{G_n\}_n \), we see that
\[ \left\| \sum_{j=J_0}^{J_0+k} |\partial_x^\sigma G_n^j \psi^j|_{\tilde{M}_{2,\delta}^3} \right\|_{L(\mathbb{R}^+)} \leq \left( \sum_{j=J_0}^{J_0+k} \left\| |\partial_x^\sigma G_n^j \psi^j|_{\tilde{M}_{2,\delta}^3} \right\|_{L(\mathbb{R}^+)} \cap S(\mathbb{R}^+) \right)^{1/\delta} + o(1). \]

By the decoupling inequality (4.10) and the above estimates, we obtain the desired estimate. \( \square \)

**Remark 4.9.** The equation (4.16) is the main improvement due to our main theorem. In the previous result [24], the improved Strichartz estimate is valid only for the diagonal case. Hence, we use a substitute by interpolating diagonal improved estimate and non-diagonal estimate in \( L^\alpha \) space. The interpolation spoils summability in \( j \), which causes a restriction on possible range of \( \alpha \).

**Proof of Proposition 4.6.** Let \( W_n^k := \sum_{j=J_0}^{J_0+k} V_n^j \), where \( J_0 \) is fixed later. Then \( W_n^k \) satisfies the integral equation
\[ W_n^k = e^{-t\partial_3^3} W_n^k(0) + \mu \int_0^t e^{-(t-s)\partial_3^3} \partial_x^\sigma W_n^k(\|W_n^k\|^{2\alpha} W_n^k + E_n^k) ds, \]
where \(-E_n^k = |W_n^k|^{2\alpha} W_n^k - \sum_{j=J_0+1}^{J_0+k} |V_n^j|^{2\alpha} V_n^j\). Applying the inhomogeneous Strichartz’ estimate [23, Proposition 2.5], we obtain
\[
\|W_n^k\|_{L^\beta(R_+) \cap S(R_+)} \leq \|e^{-i\theta t} W_n^k(0)\|_{L^\beta(R_+) \cap S(R_+)} + C\|W_n^{k+1}\|_{L^\beta(R_+) \cap S(R_+)} + C\|E_n\|_{N(R_+)},
\]
for any \(k \geq 1\) and \(n \geq N(k)\). By the interpolation inequality and the Young inequality,
\[
\|E_n\|_{N(R_+)} \leq \|\partial_x^{1/\alpha} E_n\|_{L^{p_1} L^{q_1}(R \times R_+)}^{1/2} \|E_n\|_{L^{p_2} L^{q_2}(R \times R_+)}^{1/2}
\]
and 1/\(p_1\) = 1/(L) + 2\(\alpha/p(S)\), 1/\(q_1\) = 1/q(L) + 2\(\alpha/q(S)\), \(p_2 = p(S)/(2\alpha + 1)\), and \(q_2 = q(S)/(2\alpha + 1)\). We have
\[
\|\partial_x^{1/\alpha} E_n\|_{L^{p_1} L^{q_1}(R \times R_+)} \leq C \sum_{j=J_0+1}^{J_0+k} \|V_n^j\|_{L^\beta(R_+) \cap S(R_+)} \leq C k M^{2\alpha+1}
\]
in light of (4.14). On the other hand, for any \(k\), we have \(\|E_n\|_{L^{p_1} L^{q_1}(R \times R_+)} \to 0\) as \(n \to \infty\) (see [23, Lemma 4.8]). Thus, we obtain (4.19) for any \(k \geq 1\) and \(n \geq N(k)\). Combining (4.17), (4.18), (4.19) and the continuity argument, we have that if \(\varepsilon\) is sufficiently small, then \(\|W_n^k\|_{L^\beta(R_+) \cap S(R_+)} \leq C \varepsilon\) for any \(k \geq 1\) and \(n \geq N(k)\). Combining this with (4.9), we obtain the uniform estimate (4.15).

4.3. Proof of Theorem 4.14. We finally consider analysis of \(E_2\).

Proof. By definition of \(E_2\), it is possible to choose a minimizing sequence of solutions \(\{u_n(t)\}\) so that all \(u_n(t)\) does not scatter forward in time and
\[
E_2 \leq \lim_{t \to T_{\text{max}}(u_n)} \|\partial_x^{\sigma} u_n(t)\|_{M^{\beta}_{2,\delta}} \leq E_2 + \frac{1}{n}.
\]

Hence, there exists \(t_n, t_n' \in I_{\text{max}}(u_n)\), \(t_n < t_n'\), so that
\[
\|u_n\|_{S([t_n, t_n'])} \geq n, \quad \sup_{t \in [t_n, T_{\text{max}}]} \|\partial_x^{\sigma} u_n(t)\|_{M^{\beta}_{2,\delta}} \in \left[ E_2, E_2 + \frac{2}{n} \right].
\]

Indeed, we first choose \(t_n\) so that the second property holds. Then, since \(\|u_n\|_{S([t_n, T_{\text{max}}])} = \infty\), we can choose \(t_n'\) so that the first property is true.

By time translation symmetry, we may suppose that \(t_n' = 0\). We now apply linear profile decomposition to \(u_n(0)\) to get the decomposition
\[
u_n(0) = \sum_{j=1}^{J} G_n^j \psi^j + r_n^J
\]
for \(n, J \geq 1\) with the properties (4.9), (4.11), and pairwise orthogonality of \(\{G_n^j\}\) C. By extracting subsequence and changing notations if necessary,
we may assume that for each $j$ and \( \{x_n^j\}_{n,j} = \{\log N_n^j\}_{n,j}, \{s_n^j\}_{n,j}, \{y_n^j\}_{n,j} \), we have either $x_n^j \equiv 0$, $x_n^j \to \infty$ as $n \to \infty$, or $x_n^j \to -\infty$ as $n \to \infty$. Let us define nonlinear profile $\Psi^j$ associated with $\langle \psi^j, s_n^j \rangle$ in the same way as in the proof of Theorem 1.12. We also define $V_n^j$ and $\hat{u}_n^j$ by (4.12) and (4.13), respectively.

Then, mimicking the proof of Theorem 1.12, one sees that at least one $\Psi^j$ does not scatter forward in time. We further see from decoupling inequality (4.10) and small data scattering that the number of the profiles that do not scatter is finite. Renumbering, we may suppose that $\Psi^j$ does not scatter forward in time if and only if $j \in [1, J_1]$. Here, $1 \leq J_1 < \infty$. Arguing as in [22], we see that $J_1 = 1$, $\lim_{t \to T_{\max}(\psi^j)} \| \partial_x |^{\sigma} \Psi^j(t) \|_{\tilde{M}^{\beta}_{2,\delta}} = E_2$, $\psi^j \equiv 0$ for $j \geq 2$, and $r_n^j \to 0$ as $n \to \infty$ in $|\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$. As a result,

\[
(4.20) \quad u_n(0) = G_n^1 \psi^1 + o_n(1) \quad \text{in } |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}.
\]

If $s_n^1 \to \infty$ as $n \to \infty$ then $\Psi^1(t)$ scatter forward in time, a contradiction. Because of $\|u_n\|_{S(\{u_n(0)\})} \geq n$, the same argument works for negative time direction. We see that $\Psi^1(t)$ does not scatter backward in time and that the case $s_n^1 \to -\infty$ as $n \to \infty$ is excluded. Moreover, together with $\sup_{t \in [n, T_{\max}(\psi^1)]} \| \partial_x |^{\sigma} u_n(t) \|_{\tilde{M}^{\beta}_{2,\delta}} \in [E_2, E_2 + \frac{2}{n}]$, we have

\[
\lim_{t \to T_{\min}(\psi^1)} \| \partial_x |^{\sigma} \Psi^1(t) \|_{\tilde{M}^{\beta}_{2,\delta}} = \sup_{t \in I_{\max}(\psi^1)} \| \partial_x |^{\sigma} \Psi^1(t) \|_{\tilde{M}^{\beta}_{2,\delta}} = E_2.
\]

So far, we have proven that $\Psi^1$ satisfies the first two properties of Theorem 1.14. Let us finally prove the precompactness modulo symmetry. Take an arbitrary sequence $\{\tau_n\} \subset I_{\max}(\psi^1)$. Then, we can choose $t_n \in (T_{\min}(\psi^1), \tau_n)$ so that $u_n(t) := \Psi$, $t_n^j = \tau_n$, and this $t_n$ satisfies the same assumption as above. The decomposition (4.20) reads as existence of $\psi \in |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}$, $\{N_n\} \subset \mathbb{R}_+$, and $\{y_n\} \subset \mathbb{R}$ such that

\[
\Psi^1(\tau_n) = D(N_n) T(y_n) \phi + o_n(1) \quad \text{in } |\partial_x|^{-\sigma} \tilde{M}^{\beta}_{2,\delta}.
\]

This is nothing but a sequential version of precompactness. A standard argument then upgrades this property to the continuous one. 

**Appendix A. Embedding in the generalized Morrey space**

In this appendix we mention the embedding properties of the generalized Morrey space. We first note that $M^{\beta}_{1,\infty}$ is a usual Morrey space. We easily see that $M^{\beta}_{1,\infty} = L^{\beta}$ with equal norm.

We collect the inclusion relations for the generalized Morrey space.

(i) For any $1 \leq \gamma_2 \leq \gamma_1 \leq \beta \leq \infty$ and $1 \leq \delta_1 \leq \delta_2 \leq \infty$, it holds that $M^{\beta}_{\gamma_1,\delta_1} \hookrightarrow M^{\beta}_{\gamma_2,\delta_2}$.

(ii) For any $1 \leq \beta \leq \gamma_1 \leq \gamma_2 \leq \infty$ and $1 \leq \delta_1 \leq \delta_2 \leq \infty$, it holds that $\tilde{M}^{\beta}_{\gamma_1,\delta_1} \hookrightarrow \tilde{M}^{\beta}_{\gamma_2,\delta_2}$.

(iii) $L^\beta \hookrightarrow M^\beta_{\gamma,\delta}$ holds as long as $1 \leq \gamma < \beta < \delta \leq \infty$.

(iv) $\hat{L}^\beta \hookrightarrow \tilde{M}^\beta_{\gamma,\delta}$ holds as long as $1 \leq \gamma' < \beta' < \delta \leq \infty$. 

(v) $|\partial_x|^{-\sigma} M_{\gamma_1, \delta_1}^{\beta} \lesssim M_{\gamma_2, \delta_2}^{\alpha}$ if $1/\gamma_1 - 1/\gamma_2 > 1/\beta - 1/\alpha = \sigma$ and $\delta_1 \leq \delta_2$.

The properties (i) and (ii) are trivial from the definition of the generalized Morrey space. For the proof of (iii) and (iv), see [24, Proposition A.1]. We now give the proof of (v). The Hölder inequality yields

$$
\| f \|_{L^2(\tau_k^j)} \lesssim \| |\xi|^{-\sigma} \|_{L^{2^*}(\tau_k^j)} \| |\xi|^{\alpha} f \|_{L^{2'}(\tau_k^j)} = C \| |\tau_k^j|^{-\frac{1}{2} - \frac{1}{\alpha} - \sigma} \| \| |\xi|^{\alpha} f \|_{L^{2'}(\tau_k^j)}
$$

if $1/\gamma_1 - 1/\gamma_2 - \sigma > 0$. Hence, it follows that

$$
\| f \|_{\dot{M}_{\gamma_2, \delta_2}^{\alpha}} = \| |\tau_k^j|^{-\frac{1}{2} - \frac{1}{\alpha} - \sigma} \|_{L^{2'}(\tau_k^j)} \|_{\ell^2_{k,j}} \leq C \| |\tau_k^j|^{-\frac{1}{2} - \frac{1}{\alpha} - \sigma} \| \| |\xi|^{\alpha} f \|_{L^{2'}(\tau_k^j)} \|_{\ell^2_{k,j}} \leq C \| |\partial_x|^{-\sigma} f \|_{\dot{M}_{\gamma_1, \delta_1}^{\alpha}}
$$

as long as $1/\gamma_1 - 1/\gamma_2 > \sigma = 1/\beta - 1/\alpha$ and $\delta_1 \leq \delta_2$. Hence we have (v).

Acknowledgments. S.M. is partially supported by the Sumitomo Foundation, Basic Science Research Projects No. 161145. J.S. is partially supported by JSPS, Grant-in-Aid for Young Scientists (A) 25707004.

References

[1] Béguin P. and Vargas A., Mass concentration phenomena for the $L^2$-critical nonlinear Schrödinger equation. Trans. Amer. Math. Soc. 359 (2007), 5257–5282.

[2] Benedek A. and Panzone R., The space $L^p$, with mixed norm, Duke Math. J. 28 (1961) 301–324.

[3] Bourgain J., On the restriction and multiplier problems in $\mathbb{R}^3$. Geometric aspects of functional analysis (1989–90), Lecture Notes in Math., 1469, Springer, Berlin, (1991), 179–191.

[4] Bourgain J., Some new estimates on oscillatory integrals. Essays on Fourier analysis in honor of Elias M. Stein (Princeton, NJ, 1991), Princeton Math. Ser., 42, Princeton Univ. Press (1995), 83–112.

[5] Bourgain J., Refinements of Strichartz’ inequality and applications to 2D-NLS with critical nonlinearity. Internat. Math. Res. Notices 1998 (1998), 253–283.

[6] Carles R., and Keraani S., On the role of quadratic oscillations in nonlinear Schrödinger equations II. The $L^2$-critical case. Trans. Amer. Math. Soc. 359 (2007), 33–62.

[7] Cazenave T., “Semilinear Schrödinger equations”. Courant Lecture Notes in Mathematics, 10. American Mathematical Society (2003).

[8] Christ F.M. and Weinstein M.I., Dispersion of small amplitude solutions of the generalized Korteweg-de Vries equation. J. Funct. Anal. 100 (1991) 87–109.

[9] Deift P. and Zhou X., A steepest descent method for oscillatory Riemann-Hilbert problems. Asymptotics for the MKdV equation. Ann. of Math. (2) 137 (1993), 295–368.

[10] Dodson. B., Global well-Posedness and scattering for the defocusing, mass-critical generalized KdV equation. Ann. PDE 3 (2017), Article no.5.

[11] Grünrock A., An improved local well-posedness result for the modified KdV equation. Int. Math. Res. Not. 2004 (2004), 3287–3308.

[12] Hayashi N. and Naumkin P.I., Large time asymptotics of solutions to the generalized Korteweg-de Vries equation. J. Funct. Anal. 159 (1998) 110–136.

[13] Hyakuna R. and Tsutsumi M., On existence of global solutions of Schrödinger equations with subcritical nonlinearity for $L^p$-initial data. Proc. Amer. Math. Soc. 140 (2012), 3905–3920.

[14] Kenig C.E. and Merle F., Global well-posedness, scattering and blow-up for the energy-critical, focusing, non-linear Schrödinger equation in the radial case. Invent. Math. 166 (2006), 645–675.
Kenig C.E., Ponce G. and Vega L., Oscillatory integrals and regularity of dispersive equations. Indiana Univ.math J. 40 (1991), 33–69.

Kenig C.E., Ponce G. and Vega L., Well-posedness and scattering results for the generalized Korteweg-de Vries equation via the contraction principle. Comm. Pure Appl. Math. 46 (1993), 527–620.

Kenig C.E., Ponce G. and Vega L., On the concentration of blow up solutions for the generalized KdV equation critical in $L^2$. Nonlinear wave equations (Providence, RI 1998), Contemp. Math. 263, Amer. Math. Soc., Providence, RI (2000), 131–156.

Killip R, Kwon S., Shao S. and Visan M., On the mass-critical generalized KdV equation. Discrete Contin. Dyn. Syst. 32 (2012), 191–221.

Koch H. and Marzuola J.L., Small data scattering and soliton stability in $\dot{H}^{-\frac{1}{6}}$ for the quartic KdV equation. Anal. PDE 5 (2012), 145–198.

Korteweg D. J. and de Vries G., On the change of form of long waves advancing in a rectangular canal, and on a new type of long stationary waves, Philos. Mag. 39 (1895), 422-443.

Kwon S. and Roy T., Bilinear local smoothing estimate for Airy equation. Differential Integral Equations 25 (2012), 75–83.

Masaki S., Two minimization problems on non-scattering solutions to mass-subcritical nonlinear Schrödinger equation. preprint available at arXiv:1605.09234 (2016).

Masaki S. and Segata J., On well-posedness of generalized Korteweg-de Vries equation in scale critical $\hat{L}^r$ space. Anal. and PDE. 9 (2016), 699–725.

Masaki S. and Segata J., Existence of a minimal non-scattering solution to the mass-subcritical generalized Korteweg-de Vries equation, preprint available at arXiv:1602.05331.

Merle F. and Vega L., Compactness at blow-up time for $L^2$ solutions of the critical nonlinear Schrödinger equation in 2D. Internat. Math. Res. Notices 1998 (1998), 399–425.

Moyua, A., Vargas, A., and Vega, L., Schrödinger maximal function and restriction properties of the Fourier transform. Internat. Math. Res. Notices 1996 (1996), 793–815.

Moyua, A., Vargas, A., and Vega, L., Restriction theorems and maximal operators related to oscillatory integrals in $\mathbb{R}^3$. Duke Math. J. 96 (1999), no. 3, 547–574.

Shao S., The linear profile decomposition for the Airy equation and the existence of maximizers for the Airy Strichartz inequality. Anal. PDE 2 (2009), 83–117.

Strichartz R.S., Restriction of Fourier transform to quadratic surfaces and decay of wave equation. Duke Math. J. 44 (1977), 705-714

Tao T., Scattering for the quartic generalised Korteweg-de Vries equation. J. Differential Equations 232 (2007), 623–651.

Tao T., Vargas A., and Vega L., A bilinear approach to the restriction and Kakeya conjectures. J. Amer. Math. Soc. 11 (1998), 967–1000.

Tomas P. A., A restriction theorem for the Fourier transform. Bull. Amer. Math. Soc. 81 (1975), 477–478.

Triebel H., “Theory of function spaces”. Monographs in Mathematics, 78 Birkhäuser Verlag, Basel (1983).

Vargas A. and Vega L., Global wellposedness for 1D non-linear Schrödinger equation for data with an infinite $L^2$ norm. J. Math. Pures Appl. (9) 80 (2001), 1029–1044.