Structural and Dielectric Properties of (1-x) (Al0.2La0.8TiO3) + (x) (BiZnFeO3) (x = 0.2 - 0.8) Nanocomposites

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Structural and dielectric properties of \((1-x)\) \((\text{Al}_{0.2}\text{La}_{0.8}\text{TiO}_3) + (x)\) \((\text{BiZnFeO}_3)\) \((x = 0.2 - 0.8)\) nanocomposites

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Abstract

\((1-x)\) \((\text{Al}_{0.2}\text{La}_{0.8}\text{TiO}_3) + (x)\) \((\text{BiZnFeO}_3)\) \((x = 0.2 - 0.8)\) [ALTBZFO] nanocomposites were synthesized via hydrothermal method. The X-ray diffraction patterns indicated the phase transformation from tetragonal to cubic for \(x = 0.2\) to \(0.4 - 0.8\) samples, respectively. The surface morphology showed the existence of nanospheres like structures. At 1 MHz frequency also, the dielectric constant was increased from 230 to 710 for \(x = 0.2 - 0.6\) samples, respectively. But, interestingly, \(x = 0.6\) nanocomposite exhibited the negative dielectric behavior having the dielectric constant \((\varepsilon')\) \(-58.5\) and dielectric loss \((\varepsilon'')\) \(-417\) at 8 MHz. Likewise, \(x = 0.6\) sample showed ac-electrical conductivity \((\sigma_{ac})\) \(-0.159\ S/cm\) at 6 MHz. Hence, these kinds of materials can provide high charge stored capacitor, and perfect absorber applications.

Keywords: Structural Transformation; Electron Microscopy; Negative Dielectric Constant; Diaelectricity; Impedance Spectroscopy.

1. Introduction

It was a familiar fact that the perovskite nanoceramic composites performed advanced structural, and dielectric properties [1]. Even the applications of several perovskite nanocomposites were extended to biomedical, magnetic, electromagnetic, microwave, ferroelectric, charge storage etc., [1]. Few of the above merits were seen in the case of \(\text{LaTiO}_3\) [2]. In the literature [1 - 4], \(\text{AlLaTiO}_3\) was studied extensively by several researchers for structural, optical, morphological, electrical, and dielectric properties. The results
indicated that the huge dielectric response was obtained from the compositions of AlLaTiO$_3$. Further, the nanocomposites were synthesized including the AlLaTiO$_3$, and studied for the similar properties. These materials were of ALTBT [1], ALTBFO [5, 6], LCET [7], and AELTO [2] revealing the nanorod formation, negative dielectric behavior, and defect dipole formation, respectively. This indicated a fact that the more advanced properties were achieved on mixing the barium titanate, bismuth iron oxide, and europium oxide.

Similarly, for ALTBFO nanocomposites, the zinc element was substituted within the BiFeO$_3$ system. That is, the BiZnFeO$_3$ was mixed in the ALTBZFO nanocomposite system. In general, the ZnO is a compound having variety of applications such as wide band gap semiconductor, high conducting due to high mobility, good luminescence, etc., [8]. Due to these factors, we planned to select this compound as the substituent in the composite system thereby achieving more advanced structural, morphological, dielectric, and impedance spectroscopy properties. In addition, BiFeO$_3$ showed the biosensor [9], magnetoelectric [10], optical [11], magnetic [12], multiferroic [13], enhanced electrical properties [13], and high dielectric constant [14]. Further, it was also noticed in the literature that the BiZnFeO$_3$ provided photovoltaic [15] and progressed magnetic properties [16]. In case of Zn-doped BiFeO$_3$, the leakage, and high polarization behavior was observed predominantly [17]. Similarly, there was a significant influence on optical, electrical, and magnetic properties upon doping the zinc in the bismuth iron oxide perovskite system [18 - 20]. Uniyal et al. [21], reported the multiferroic nature of ZnFe$_2$O$_4$–BiFeO$_3$ nanocomposites. Hence, the ALTBZFO nanocomposites were prepared using hydrothermal method owing to the benefits such as low operating temperature, less time & power consuming, well crystallinity, low particle size, etc., [1].
2. Materials and Methods

For the synthesis of \((1-x)\) \((\text{Al}_{0.2}\text{La}_{0.8}\text{TiO}_3) + (x)\) \((\text{BiZnFeO}_3)\) \((x = 0.2 - 0.8)\) [ALTBZFO] nanocomposites, the hydrothermal method was preferred. As a part of this, the raw materials were considered to be \(\text{Al(NO}_3\text{)}_3\cdot9\text{H}_2\text{O}\) (99.8 % purity, Sigma-Aldrich), \(\text{La(NO}_3\text{)}_3\cdot6\text{H}_2\text{O}\) (99.8 % purity, Sigma-Aldrich), \(\text{Bi(NO}_3\text{)}_3\cdot5\text{H}_2\text{O}\) (99.8 % purity, Sigma-Aldrich), \(\text{Fe(NO}_3\text{)}_3\cdot9\text{H}_2\text{O}\) (99.9 % purity, Sigma-Aldrich), \(\text{Zn(NO}_3\text{)}_2\cdot6\text{H}_2\text{O}\) (99.9 % purity, Sigma-Aldrich), and \(\text{TiO}_2\) (99.9 % purity, Sigma-Aldrich). These precursors were taken in a stoichiometric ratio and kept in a cleaned glass beaker. Further, the mixture was dissolved in distilled water, and stirred for half an hour at room temperature by maintaining stirring rate of 450 rpm. As a result of this, a delicious solution was formed, and it was kept in well cleaned Teflon bowl of 300 ml capacity. This bowl was again placed in a stainless-steel autoclave. Then, the autoclave was kept in a hot air oven, and further the reaction was taken place at 150°C/6 hr. After completion of reaction, the autoclave was quenched to room temperature. Then, the Teflon bowl was removed from the autoclave to extract the nanocomposite solution. This solution was washed for 10 times, and further dried for 2 hr., at 60°C. Thus, the final nanocomposite product was obtained for \(x = 0.2 - 0.8\) samples (Fig.1). In the next step, the nanocomposites were subjected to various analysis works like X-ray diffractometer (XRD, Bruker, \(\text{CuK}_\alpha\), \(\lambda = 0.15406\) nm), transmission electron microscope (TEM: Model Tecnai G20, FEI, USA), field-emission scanning electron microscope (FE-SEM, Ultra 55, Carl Zeiss), and LCR controller (HIOKI 3532-50, 42 Hz–8 MHz) to study the structural, morphological, and dielectric properties, respectively.
3. Results and discussion

The X-ray diffraction patterns of ALTBZF [(1-x) (Al_{0.2}La_{0.8}TiO_3) + x (BiZnFeO_3)] (x = 0.2 - 0.8) nanocomposites were recorded by considering the two-theta angle on X-axis and the corresponding intensity all reflection planes on Y-axis. From Fig. 2, it was observed that all the nanocomposites showed well crystalline nature. It was also seen that the maximum intensity of x = 0.2 sample showed at 2θ = 25.463°, while the x = 0.4 - 0.8 samples revealed the maximum intensity at 2θ = 31.421°. This established a fact that there was a phase transition from x = 0.2 to 0.4 - 0.8. All the diffraction positions of x = 0.2 - 0.8 composites were compared with standard JCPDS. Then it was observed that the x = 0.2 sample showed the reflection planes comparable to the JCPDS: 21 – 1272, while x = 0.4 – 0.8 nanocomposites contain the reflections comparable to the JCPDS: 86 – 0368. This comparison confirmed a fact that the x = 0.2 sample achieved the tetragonal structure, whereas the x = 0.4 – 0.8 samples showed the cubic phases. In addition, the secondary phases related to the BiFeO_3 (indicated by *) and La_2TiO_3 (indicated by +) were noticed. Usually, this kind of secondary phase formation was attributed to the vast difference of ionic radii of cations present in the composite system. That is, La^{3+}: 0.136 nm, Bi^{3+}: 0.096 nm, Al^{3+}: 0.039
nm, Ti$^{4+}$: 0.068 nm, Fe$^{3+}$: 0.0645 nm, and Zn$^{2+}$: 0.060 nm [22]. Herein, the bismuth, and lanthanum cations have larger ionic radii as compared with the rest of the cations. Therefore, bismuth, and lanthanum cations can form their associated compounds (BiFeO$_3$ and La$_2$TiO$_5$). The high crystallinity was also identified for all diffraction planes. Further, the average crystallite diameter was calculated using the Scherrer relation [23]: $D_a = 0.9\lambda/\beta\cos\theta$, where $\lambda$ indicates the wavelength of X-rays used (0.15418 nm), $\beta$ is full width half maxima (FWHM) and $\theta$ is the angle of diffraction. The results indicated that the $D_a$ was found to be increasing from 75.3 to 160.3 nm for $x = 0.2 – 0.6$, respectively, while for $x = 0.8$, it was noticed to be 107.4 nm (decreased). This kind of manner was attributed to the decrease of FWHM (from 0.443 – 0.104 radian) for $x = 0.2 – 0.6$ samples, and further it was increasing to 0.109 radian. This revealed a fact that there exists an inversely proportional relationship with between FWHM and $D_a$. Afterwards, the lattice parameters (a, b & c) were determined using the relations [5, 24, 25]: $1/d^2 = (h^2+k^2)/a^2 + l^2/c^2$ (for tetragonal) and $1/d^2 = (h^2+k^2+l^2)/a^2$ (for cubic). The obtained results indicated that for $x = 0.2$ sample, $a = b = 3.774$ Å & $c = 9.395$ Å (Table. 1). On the other hand, for $x = 0.4 – 0.8$ samples expressed the decreasing trend from 12.761 to 12.714 Å, respectively. These results clearly indicated that the structural transformation was occurred from tetragonal to cubic for $x = 0.2$ to 0.4/0.6/0.8 samples. Similarly, the unit cell volume (V) was found to be 133.81 Å$^3$ for $x = 0.2$ while the same parameter was decreased from 2078.1 to 2055.2 Å$^3$, for $x = 0.4 – 0.8$, respectively. Furthermore, the X-ray density values were calculated using a formula: $\rho_x$ (g/c.c.) = $Z$*C.M.W./(N*$V$) (for tetragonal $Z = 4$ & for cubic $Z = 1$), where $Z$ = effective number of atoms per unit cell, C.M.W is the composite’s molecular weight, N is the Avogadro’s number (6.023x10$^{23}$), and $V$ is the unit cell volume [5]. The $\rho_x$ of $x = 0.2$ was about 12.187 g/c.c., while the similar parameter was increased from 0.223 to 0.279 g/c.c., for $x = 0.4 – 0.8$ samples, respectively. This kind of behavior was attributed to the increasing trend of
composite molecular weight (C.M.W.) from 245.55 to 345.04 g/mole. Moreover, the specific surface area (S) was noted to be very low for x = 0.2, whereas for x = 0.4 – 0.8 samples, it was found to be changing from 149.7 to 248.3 m$^2$/g. The high ‘S’ values usually indicate the advanced electrical properties.

Table.1 XRD parameters of ALTBZF nanocomposites.

| x  | 0.2  | 0.4  | 0.6  | 0.8  |
|----|------|------|------|------|
| $D_a$ (nm) | 75.3 | 108.5 | 160.3 | 107.4 |
| FWHM (radian) | 0.443 | 0.111 | 0.104 | 0.109 |
| $a = b$ (Å) | 3.774 | 12.761 | 12.747 | 12.714 |
| $c$ (Å) | 9.395 | 12.761 | 12.747 | 12.714 |
| $V$ (Å$^3$) | 133.81 | 207.81 | 2071.2 | 2055.2 |
| C.M.W (g/mole) | 245.55 | 278.71 | 311.88 | 345.04 |
| $\rho_x$ (g/c.c.) | 12.187 | 0.223 | 0.250 | 0.279 |
| S (m$^2$/g) | 6.538 | 248.3 | 149.7 | 200.4 |

Fig.2 XRD patterns of ALTBZF nanocomposites
The surface morphology of ALTBZF nanocomposites was analyzed using the field emission scanning electron microscopy (FESEM), and transmission electron microscopy (TEM) pictures. The FESEM pictures were shown in Fig.3 at 100 nm scale. In Fig.3, it was observed that the larger numbers of cotton spheres like grains were formed for $x = 0.2$. Moreover, few nanofibers like grains were also deposited in the microstructure. It was also noticed that the distribution of cotton spheres was almost homogeneous. For $x = 0.4$ samples, the bigger sized and well-defined cotton spheres were seen. But for $x = 0.6$ sample exhibited the clustered like grains possessing the asymmetrical cotton spheres. On the other hand, the $x = 0.8$ sample revealed the existence of cotton spheres in limited number covered by smaller nanofibers. It was understood that with an increase of BiZnFeO$_3$ in the composite system, the morphology was changed from larger numbers of cotton spheres to smaller nanofibers including limited numbers of cotton spheres. In the same way, the TEM pictures were recorded as provided in Fig.4. From TEM pictures, we noticed that $x = 0.2 – 0.6$ samples exhibited the formation of symmetrical nanospheres including small nanofibers. On the other hand, $x = 0.8$ sample showed the presence of nanospheres including larger nanofibers as compared with $x = 0.2 – 0.6$ samples. The similar kind of morphology was observed in the FESEM pictures also. Further, the average particle size was measured and found to be altering from 115 to 166 nm for all the nanocomposites. Even the selected area electron diffraction patterns (SAED) (Fig.5) were recorded and showed the presence of concentric circular rings with good intensity. These diffraction rings agreed with X-ray diffraction patterns.
Fig. 3 FESEM pictures of ALTBZF nanocomposites
Fig. 4 TEM pictures of ALTBZF nanocomposites
Fig. 5 SAED patterns of ALTBZF nanocomposites
The dielectric behavior of ALTBZFO nanocomposites was explained using the dielectric constant ($\varepsilon'$), and dielectric loss ($\varepsilon''$) versus frequency (log $\omega$) plots. In Fig.6 ($\varepsilon'$ - log $\omega$), it was seen that the $\varepsilon'$ was very high for $x = 0.2 - 0.8$ samples. Usually, this kind of behavior can be attributed to the space charge species. That is, the charges will be piled up at the grain boundary interface at low log $\omega$. Thus, the polarization can be developed progressively, and it can be treated as Maxwell-Wagner’s polarization [26]. It leads achieving the high dielectric constant. It was seen from the $\varepsilon'$ - log $\omega$ plots that the $\varepsilon'$ was found to be increasing from 3727 to 14166 for $x = 0.2 - 0.6$ samples (Table.2), respectively at 1 kHz. But for $x = 0.8$ sample, it was decreased to 5081. Similarly, the $\varepsilon''$ (in Fig.6 ($\varepsilon''$ - log $\omega$)) also showed the increasing trend from 566 to 15340 for $x = 0.2 - 0.6$ samples, respectively. But for $y = 0.8$, it was reduced to 1509. At 1 MHz frequency also, the similar increasing trend was occurred from 230 to 710 for $x = 0.2 - 0.6$ samples, respectively. For $x = 0.8$, it was decreased to 252. From 6 to 8 MHz also, the similar behavior was observed for all nanocomposites except for $x = 0.6$. For all the samples, the dielectric relaxations were noticed in the dielectric loss versus frequency plots at log $\omega = 5$ to 6.35. These relaxations were formed normally, due to the accumulation of charges at the interface of sample. Moreover, interestingly, $x = 0.6$ sample showed negative dielectric behavior having the $\varepsilon' \sim -58.5$ $\varepsilon'' \sim -417$ (Table.2). This indicated a fact that for $x = 0.6$ sample, the transition was happened from positive to negative dielectric behavior at 6 MHz frequency. In general, it was known that the positive dielectric behavior is a usual property while the negative dielectric behavior is an unusual property. This kind of property was earlier observed in the case of perovskite materials [27, 28] as well as the metamaterials [2, 7, 29 - 33]. However, the present nanocomposites are not the metamaterials. But nevertheless, the negative dielectric behavior can be obtained due to the geometrical effect thereby providing the negative polarization. That is, for the provided electric field frequency, the charge carriers will be aligned in the
opposite direction of the applied electric field. Therefore, the polarization will be developed having negative sign. This can lead to the formation of negative polarization. Consequently, the negative dielectric constant, and dielectric loss will be achieved. This behavior can be considered as dielectricity nature which is like the diamagnetism [2]. That means, in case of diamagnetic nature of a magnetic specimen, the magnetization will be developed opposite to the direction of applied magnetic field. Likewise, the in case of dielectricity, the polarization will be developed opposite to the direction of applied electric field frequency. The similar observations were evidenced in the case of $x = 0.6$ nanocomposite during $6–8$ MHz frequency. From these results, it was understood that the nanocomposites ($x = 0.2, 0.4, \text{ and } 0.8$) exhibiting the high dielectric constant can be suited for high charge storage capacitor applications [2]. Likewise, the $x = 0.6$ sample offering negative dielectric behavior can provide applications in waveguides, filters, antenna, electromagnetic cloaking devices, and perfect microwave absorbers [2]. Further, the ac-electrical conductivity was calculated using a relation $\sigma_{ac}: \varepsilon_0 \varepsilon'' \omega$, where $\varepsilon_0$ is the permittivity of free space, and the other symbols have their usual meaning [34]. To explain the electrical conductivity as a function of composite and frequency, the log $\sigma_{ac}$ versus log $\omega$ plots (Fig.7) were drawn. It was clear that at low log $\omega$ value, the $\sigma_{ac}$ was very small while going to the high log $\omega$ values, the conductivity was increased progressively. This was happened for $y = 0.2, 0.4, \text{ and } 0.8$ samples. But for $x = 0.6$ sample, the ac-electrical conductivity was noticed to be negative from 6 to 8 MHz frequency. At 1 MHz, the $\sigma_{ac}$ was found to be varying from 0.0142 to 0.0976 S/cm for all nanocomposites. This indicated that the high conductivity was noticed for all the samples at room temperature (at 1 MHz). For $x = 0.6$ sample, the $\sigma_{ac}$ showed -0.159 S/cm at 6 MHz (Table. 3). This manner was inferred due to the negative dielectric loss of the same sample.
Fig. 6 Frequency dependence of dielectric properties of ALTBZF nanocomposites

Table 2. Dielectric parameters of ALTBZF nanostructures at different frequencies

| x  | 0.2   | 0.4   | 0.6   | 0.8   |
|----|-------|-------|-------|-------|
| frequency | $\varepsilon'$ | $\varepsilon''$ | $\varepsilon'$ | $\varepsilon''$ | $\varepsilon'$ | $\varepsilon''$ | $\varepsilon'$ | $\varepsilon''$ |
| 1 kHz | 3727  | 566   | 4700  | 755   | 14166 | 15340 | 5081  | 1509   |
| 10 kHz| 3123  | 646   | 3457  | 1226  | 8292  | 3715  | 3424  | 1415   |
| frequency  | 1 MHz  | 2 MHz  | 3 MHz  | 4 MHz  | 5 MHz  | 6 MHz  | 7 MHz  | 8 MHz  |
|-----------|--------|--------|--------|--------|--------|--------|--------|--------|
|           | 1771   | 280    | 98     | 82     | 74     | 68     | 64     | 58     |
|           | 1154   | 132    | 209    | 174    | 152    | 142    | 139    | 142    |
|           | 1013   | 172    | 112    | 102    | 95     | 92     | 89     | 87     |
|           | 1198   | 254    | 119    | 119    | 110    | 107    | 108    | 114    |
|           | 5565   | 69     | 128    | 69     | 31     | 60     | -2.4   | -58.5  |
|           | 2864   | 603    | 751    | 60     | 513    | -457   | -432   | -417   |
|           | 975    | 53     | 103    | 49     | 49     | 48     | 46     | 45     |
|           | 1237   | 86     | 105    | 75     | 75     | 70     | 67     | 68     |
|           |        |        |        |        |        |        |        |        |

Table 3. AC-electrical conductivity of ALTBZF nanostructures at different frequencies

| x   | 0.2   | 0.4   | 0.6   | 0.8   |
|-----|-------|-------|-------|-------|
| freq | 1 MHz | 6 MHz | 7 MHz | 8 MHz |
| σac (S/cm) |
| 0.0265 | 0.0153 | 0.0976 | 0.0142 |
|   -   |     -  |    -   |    -   |
|   -   |     -  |    -   |    -   |
|   -   |     -  |    -   |    -   |

Fig. 7 AC-electrical conductivity versus log ω plots of ALTBZF nanocomposites

To analyze the microstructure, electrical conduction, and polarization mechanisms, the impedance analysis was carried out. For this, the complex impedance \( Z^* = Z' - jZ'' \), wherein the \( Z' \) is the real, and \( Z'' \) is the imaginary part. In Fig. 8, the variation of \( Z' \), and \( Z'' \)
was provided as a function of log ω. The $Z'$, and $Z''$ values are very high at low log ω values. This behavior was obtained owing to the space charge effect. That is, the charge carriers will be accumulated at the interface, and therefore, this trend can lead to the high magnitude of $Z'$, and $Z''$ values. Moreover, at high log ω values, the $Z'$, and $Z''$ parameters were decreased to small values. This can indicate a fact that the effect of charge species was decreased at high frequencies. The Cole-Cole plots ($Z'$ versus $Z''$) were shown in Fig.9. For all the samples, the Cole-Cole plots showed the formation of single, and partial semicircular arcs. In general, the partial arcs will be formed due to the long-range motion of the charge carriers [34]. But for x = 0.6, the reverse nature of arc was observed at high frequencies. This manner was occurred due to the dielectricity behavior at frequencies. The dielectric modulus spectrum was analyzed using $M' – \log \omega$, $M'' – \log \omega$, and $M'$ versus $M''$. In the $M' – \log \omega$, and $M'' – \log \omega$ plots (Fig.10) [$M' = \varepsilon'/(\varepsilon'^2+\varepsilon''^2)$ (real part), and $M'' = \varepsilon''/(\varepsilon'^2+\varepsilon''^2)$ (imaginary part)], it was found that the $M'$, and $M''$ values showed almost zero value at low log ω. This behavior was attributed to the lack of control of mobility of charges and not having the restoring force to the charges to attain their original position [2]. For further increase of frequency, $M'$, and $M''$ values were increased sharply. This was attributed to the mobility of charges for short distances [2]. There were no considerable relaxations in the $M' – \log \omega$, $M'' – \log \omega$ plots. The $M'$ versus $M''$ plots (Fig.11) ensured that the single, and partial semicircular arcs were seen for x = 0.2 – 0.8 nanocomposites. At high frequencies, the x = 0.6 sample showed the complete reverse nature of arc indicating the dielectricity behavior. The similar observations were noticed in the literature [2, 7, 27, 28]. The partial arcs were formed due to the partial relaxation strength.
Fig. 8 $Z'$- log $\omega$ and $Z''$- log $\omega$ plots of ALTBZF nanocomposites
Fig. 9 $Z'$ versus $Z''$ plots of ALTBF nanocomposites
Fig. 10 $M'$- log $\omega$ and $M''$- log $\omega$ plots of ALTBZF nanocomposites

Fig. 11 $M'$ versus $M''$ plots of ALTBZF nanocomposites
4. Conclusions

The ALTBZFO nanocomposites were prepared using the low temperature hydrothermal method. The XRD analysis indicated the structural transformation from tetragonal to cubic for $y = 0.2$ to $0.4 - 0.8$ samples, respectively. The lattice constants for $x = 0.2$ sample were noted to be $a = b = 3.774 \text{ Å} \& c = 9.395 \text{ Å}$. But, for $x = 0.4 - 0.8$ samples, these lattice constants were decreased from $12.761$ to $12.714 \text{ Å}$, respectively. The FESEM, and TEM pictures revealed the generation of nanospheres, and few nanorods in the morphology. The high dielectric constant was obtained for $x = 0.2 - 0.8$ samples at $1 \text{ MHz}$ frequency. Hence, the present nanocomposites were suitable for charge stored capacitor, and perfect absorber applications. Moreover, interestingly, $x = 0.6$ nanocomposite showed the negative dielectric behavior from $6$ to $8 \text{ MHz}$. In addition, the high ac-electrical conductivity was noticed for $x = 0.2 - 0.8$ samples at $1 \text{ MHz}$. The impedance, and dielectric modulus analysis evidenced the microstructure, and conduction mechanism (behavior of charge carriers).

Conflicts of Interest: The authors declare that we have no conflicts of interest.

Data availability statement: The data will be made immediately available based on the request.

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Figure 1

Schematic representation of hydrothermal synthesis and sample preparation
Figure 2

XRD patterns of ALTBZF nanocomposites
Figure 3

FESEM pictures of ALTBZF nanocomposites
Figure 4

TEM pictures of ALTBZF nanocomposites
Figure 5

SAED patterns of ALTBZF nanocomposites
Figure 6

Frequency dependence of dielectric properties of ALTBZF nanocomposites
Figure 7

AC- electrical conductivity versus log ω plots of ALTBZF nanocomposites
Figure 8

$Z'$- log $\omega$ and $Z''$- log $\omega$ plots of ALTBZF nanocomposites
Figure 9

$Z'$ versus $Z''$ plots of ALTBZF nanocomposites
Figure 10

$M'$- log $\omega$ and $M''$- log $\omega$ plots of ALTBZF nanocomposites
Figure 11

$M'$ versus $M''$ plots of ALTBZF nanocomposites