Outside the Lyman-break box: detecting Lyman continuum emitters at $3.5 < z < 5.1$ with CLAUDS

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ABSTRACT

Identifying non-contaminated sample of high-redshift galaxies with escaping Lyman continuum (LyC) flux is important for understanding the sources and evolution of cosmic reionization. We present CLAUDS (CFHT Large Area $u$-band deep survey) $u$-band photometry of the COSMOS field to probe LyC radiation from spectroscopically confirmed galaxies at $z \geq 3.5$ and outside the standard Lyman-break galaxy colour-selection expectations. Complementary to the CLAUDS data, we use Subaru multifilter photometry, Hubble Space Telescope (HST) multifilter imaging, and the spectroscopic surveys D10K, VUDS, and 3D-HST. We present a sample of Lyman continuum galaxy (LCG) candidates in the redshift range $3.5 \lesssim z \lesssim 5.1$. Here, we introduce 5 LCG candidates, where two are flagged quality 1 and three quality 2. The estimated $f_{\text{esc}}$ for quality 1 candidates are in the range $\sim 5 - 73$ per cent and $\sim 30 - 93$ per cent. These estimates are based on our derived parameters from individual galaxies as inputs to a range of BPASS models as well as mean intergalactic medium (IGM) and maximal intergalactic and circumgalactic media (IGM+CGM) transmission. We conclude that our search for LCGs is most likely biased to lines of sight with low H I densities or free from Lyman limit systems. Our two best LCG candidates have EW (Ly$\alpha$) $\leq 50$ Å and we find no correlation or anticorrelation between EW (Ly$\alpha$), $f_{\text{esc}}$, and $R_{\text{obs}}$, the ratio of ionizing to non-ionizing observed flux in the measured passbands. Stacking candidates without solid LyC detections ($S/N < 3$) results in an estimated $f_{\text{esc}}$ from galaxies not greater than 1 per cent.

Key words: galaxies: high-redshift – galaxies: photometry – dark ages, reionization, first stars.

1 INTRODUCTION

The Epoch of Reionization (EoR) is a critical period where structures in the Universe such as stars, galaxies, quasars, and active galactic nuclei (QSO/AGNs), started to form and evolve. The EoR can be considered a transition period from a neutral and opaque Universe towards the mostly transparent and ionized Universe that we observe today. This stage of the Universe’s history took place between $6 \lesssim z < 15$, where $z \sim 6$ is estimated from observations of the Gunn–Peterson trough in the spectra of QSOs (Fan et al. 2002, 2006; Becker, Bolton & Lidz 2015; Eilers, Davies & Hennawi 2018; Bosman et al. 2018). Similar results are obtained by searching for a drop in the fraction of the Ly$\alpha$-emitting galaxies (LAEs), which predicts the end point of reionization in the range between $5.7 < z < 7$ (Kashikawa et al. 2006; Ouchi et al. 2010; Ota et al. 2010; Faisst et al. 2014; Konno et al. 2018; Mason et al. 2018). There are many indirect observational constraints that point to $z \sim 15$ as the most likely beginning of the EoR (Greig & Mesinger 2017; Bowman et al. 2018; Planck Collaboration et al. 2018).

Understanding the nature of the sources that emit Lyman continuum radiation (LyC) and reionized the Universe is one of the most persistent questions in modern astronomy. LyC is UV radiation with energy $E \geq 13.6$eV or $\lambda \leq 912$ Å that is able to ionize hydrogen. The first LyC photons are believed to have been emitted by massive objects like metal-free Population III stars (Bromm, Coppi & Larson 2002; Wyithe & Cen 2007; Ahn et al. 2012; Susa, Hasegawa & Tominaga 2014). In later stages of the EoR, LyC radiation is thought to be primarily produced by O and B stars in young star-forming galaxies (SFGs) as well as AGN (Madau, Haardt & Rees 1999; Madau & Haardt 2015). An additional complexity is that the
contribution to the ultraviolet background (UVB) from different sources is not constant, rather it changes as different populations of the objects evolve (Wyithe & Bolton 2011; Becker & Bolton 2013; Kakuchi et al. 2018). Although great progress has been made in the last 20 yr, the question of which sources are responsible for reionization of the intergalactic medium (IGM) remains open.

Clarifying the relative contribution of the various sources of ionizing radiation, particularly the relative roles of SFGs and AGN, is still under debate (Becker et al. 2015; Madau & Haardt 2015). Various studies supported by observations indicate that the population of AGN are not sufficient to ionize the IGM at z > 3 (Haardt & Madau 1996; Cowie, Barger & Trouille 2009; Fontanot, Cristiani & Vanzella 2012; Grissom, Ballantyne & Wise 2014; Trebitsch et al. 2018; Kulkarni, Worseck & Hennawi 2019) and sufficiently contribute to the observed global UVB, as the contribution of AGN to the UVB peaks at z ∼ 2 (Cowie & Barger 2008). However, some recent research points toward low-luminosity AGN at high redshifts as a possible main driver of reionization at its early stages (Giallongo et al. 2012, 2015; Grazian et al. 2018).

Although AGNs may contribute to the UVB, currently there is widespread acceptance that the major producers of LyC radiation are young SFGs. Therefore, galaxies are considered the most likely sources responsible for driving the reionization of the Universe. Due to the fact that LyC radiation is in the far-UV part of the spectrum, we are limited to observing galaxies at roughly 3 ≤ z ≤ 4.5 using ground-based telescopes. At z > 3, the redshifted LyC flux falls in the optical part of the spectrum and can be observed from the ground. The z ≤ 4.5 limit comes from the fact that number density of systems containing neutral hydrogen rapidly increases at z > 4 (Sargent, Steidel & Boksenberg 1989; Inoue & Iwata 2008) and the chances of detecting LyC flux decreases to below 20 per cent (Inoue & Iwata 2008).

In the last 20 yr, many observational (spectroscopic and photometric) efforts have been made to directly detect LyC flux from z ∼ 2.5–4.5 galaxies to measure their escape fraction of ionizing flux (fesc) and to define the population of galaxies contributing to reionization (Leitherer et al. 1995; Steidel, Pettini & Adelberger 2001; Inoue et al. 2005; Iwata et al. 2009; Nestor et al. 2011; Vanzella et al. 2012; Grazian et al. 2012; Siana et al. 2015; Vasei et al. 2016; Rutkowski et al. 2017; Marchi et al. 2017, 2018; Steidel et al. 2018; Fletcher et al. 2019; Bassett et al. 2019). Until now, only a few have been successful in producing spectroscopic confirmations: Ion2 (Vanzella et al. 2015; de Barros et al. 2016; Vanzella et al. 2016), Q1549–C25 (Shapley et al. 2016), Ion3 (Vanzella et al. 2018), as well recent ones (Nakajima et al. 2019; Steidel et al. 2018). It is crucial to our understanding of the EoR and structure formation to identify a larger sample; however from previous results it can be concluded that developing efficient selection criteria for LyC emitting galaxies at 3 < z < 4 remains a challenging task.

To date, galaxies examined for escaping LyC photons have been mostly selected using a variation of the Lyman break technique introduced by Steidel et al. (1996) or through narrow-band selection of the LAEs (e.g. Cowie & Hu 1998; Iwata et al. 2009; Nestor et al. 2011). However it has been recognized that with Lyman-break galaxy (LBG) selection, a significant number of high-redshift galaxies can escape the selection criteria (Steidel et al. 1999; Le Fèvre et al. 2005), as the technique is developed to be efficient but not comprehensive. The LBG colour criteria enclose an efficient selection region in colour–colour space where z ∼ 3–4 galaxies reside based on the assumption of zero escaping LyC flux. Cooke et al. (2014) applied various levels of LyC flux to LBG composite spectra and found that there is a notably large fraction (∼ 32 per cent) of z ∼ 3–4 SFGs that reside outside the standard LBG colour-selection region. Their colours are consistent with those of spectroscopically confirmed galaxies of Le Fèvre et al. (2005) randomly chosen in a ‘blind’ magnitude-limited survey (i.e. no colour selection) and galaxies selected via deep medium-band infrared photometry from the ZFOURGE survey (Straatman et al. 2016), complemented by ∼ 30 band photometry in the COSMOS, CDFS, and UDS legacy fields. The high-redshift galaxies that fall outside the LBG selection box are estimated to have medium to high fesc values and these newly identified galaxies has been termed Lyman continuum galaxies (LCGs; Cooke et al. 2014).

Complicating all selection methods for galaxies emitting LyC flux is the possibility of foreground contamination from low-redshift objects. The probability of contamination by foreground objects increases at higher redshift. Due to this kind of contamination, non-ionizing emission from a foreground galaxy can easily be mistaken for LyC radiation. Research on the estimated probability of such contamination indicates that a non-negligible fraction (∼ 7–13 per cent) of LyC candidates are contaminated by foreground objects (Siana et al. 2007; Vanzella et al. 2010a; Mostardi et al. 2013; Cooke et al. 2014). There are two ways to check high-redshift LyC candidates for contamination. The first is to use deep spectroscopy of LCG candidates to search for low-redshift galaxy features and the second is using high-spatial resolution space-based imaging i.e. Hubble Space Telescope (HST), since the contamination rate is proportional to the point spread function (Vanzella et al. 2012; Siana et al. 2015).

Another way to overcome the difficulties of studying high-redshift SFGs is to study their proxies in the local Universe. However, complications while studying low-redshift starburst galaxies arise from the fact that they are mostly opaque to the ionizing radiation that is generated in them (Grimes et al. 2009). Promising LCG counterparts at z < 1 are the Green Pea galaxies (GP), introduced by Cardamone et al. (2009), because they are compact objects with intense star formation rates. The high [O II]/[O I] ratios, high densities and possible presence of shocks found in some GPs indicate that they may be leaking ionizing photons into the IGM (Jaskot & Oey 2013; Nakajima & Ouchi 2014; Faist 2016). Research by Izotov et al. (2016, 2018b) reports the detection of escaping LyC radiation from four compact SFGs1 with fesc ∼ 6–13 per cent. More recently, the same group reported the highest fesc = 46 ± 2 per cent detection to date in low-redshift compact SFGs (Izotov et al. 2018a). These recent observational results indicate that all compact SFGs with reported LyC detections in Izotov et al. (2016, 2018a) share the same properties.

Detailed spectral analysis reveals that for all compact star forming LyC leakers the equivalent width (EW) of Lyα increases with increasing LyC escape fraction. Lyα lines in LCG leakers are also found to be double peaked. A decrease in the peak separation between the red and blue peaks is also found to correlate with an increase in fesc (Verhamme et al. 2015, 2017), Kakiichi & Gronke (2019), and Kim et al. (2019) and on average Lyα escape fraction correlates with LyC escape fraction, indicating the link between the two escape processes. All of these results appear promising in terms of

\[ f_{\text{esc}} = \frac{N_{\text{LyC}}}{N_{\text{LyC, total}}} \]

1Compact SFGs include GPs and luminous compact galaxies. Galaxies in the z ∼ 0.0–0.6 range with HÎ± line luminosity L(HÎ±) ≥ 10^{40.5} erg s^{-1} are named Luminous Compact Galaxies. General characteristics of the compact SFGs are strong emission lines in the optical part of the spectrum that are coming from H II regions, produced by ionizing radiation from O-stars (Izotov, Guseva & Thuan 2011).
of finding reliable selection criteria for high-redshift counterparts of LCGs. But first, all of these correlations need to be statistically verified from larger samples of the LyC leaking galaxies at redshifts beyond \( z \sim 3 \).

These galaxies at \( z \sim 3-5 \) are recognized as lower redshift counterparts of the galaxies responsible for reionization of the Universe for which it is impossible to directly detect ionizing LyC radiation. This is why it is crucial to our understanding of the EoR that large samples of LCGs are identified to measure ionizing LyC flux directly and to provide a sample for calibrating indirect indicators (colour, \( \text{Ly} \alpha \) line properties, \([\text{O} \text{ iii}]/[\text{O} \text{ ii}]\) ratio, etc.) that will point to the leakage of ionizing LyC radiation from \( z \sim 6 \) galaxies into IGM. However, currently developing efficient selection criteria for LCGs at \( z \sim 3-5 \) remains a challenging task.

In this work, we present results from our search for LyC ionizing radiation from \( 3.5 < z < 5.5 \) LCGs using CLAUDS photometry with the aim to test the hypothesis that these galaxies reside outside the standard LBG selection box or LBG selection expectations. In Section 2, we explain the sample selection, and in Section 4 we describe the method for LyC flux measurements. We present subsample analyses in Section 3, discuss our results in Section 5, and summarize our findings in Section 6.

2 DATA AND SAMPLE SELECTION

In order to detect LyC radiation and estimate \( f_{\text{esc}} \) from galaxies at \( z \geq 3.5 \) to study their contribution to the global budget of ionizing photons, we select and analyse candidates from the well-characterized COSMOS field (Scoville et al. 2007a). The motivations for choosing the COSMOS field are the following:

(i) The availability of spectroscopic surveys: the DEIMOS 10k spectroscopic survey (hereafter D10K; Hasinger et al. 2018), VIMOS Ultra Deep Survey (VUDS; Le Fèvre et al. 2015; Tasca et al. 2017) and as well as 3D-HST grism spectroscopy (Brammer et al. 2012; Momcheva et al. 2016).

(ii) The space-based high-resolution imaging coverage of the whole field in at least one \( \text{HST} \) filter, ACS F814W (Koekemoer et al. 2007; Scoville et al. 2007b; Massey et al. 2010), and over large areas in F125W, F160W, F140W, F606W, F336W, and F435W.

(iii) Access to the CFHT Large Area \( u\)-band deep survey (CLAUDS; Sawicki et al. 2019).

(iv) The ultradeep photometry from Hyper Suprime-Cam (HSC) Subaru Strategic Program public data release 1 (PDR1; Aihara et al. 2018a, b).

We note that to ensure the CLAUDS \( u\)-band detections are attributed correctly to pure LyC flux, it is necessary to have spectroscopic redshifts. The spectroscopic redshifts of our sample were taken from the literature; we acknowledge that getting these from the literature introduces biases in galaxy selection methods. The goal of this paper is to start with a conservatively selected spectroscopic parent sample and inspect the \( u\)-band images for LyC flux. The details of the sample selection are described below.

2.1 Photometric and spectroscopic data

The photometric data we use to probe LyC flux in our sample are provided by the CLAUDS survey whose astrometry is matched to the \( \text{Subaru} \) HSC DR1 data. The minimum depth of the CLAUDS imaging in the COSMOS field is \( \Delta m_{\text{AB}} \sim 27.2 \) AB (5\( \sigma \) in 2 arcsec diameter aperture; Sawicki et al. 2019) and it is comparable in depth with the HSC-Deep program.

![Figure 1](https://image.pollux.astro.ox.ac.uk/4918/4918351) The composite LBG spectrum from (Shapley et al. 2003) reddshifted at \( z = 3.42 \) (black) and transmission curves of the filters used in this work: CLAUDS \( u \) (blue curve) and \( \text{Subaru} \) HSC \( g, r, \) and \( i \) (green, orange, and red curves, respectively). The composite LBG spectrum is overlaid and illustrates that the \( u \) filter probe the LyC flux and \( g, r, i \) filters probe the regions of non-ionizing UV flux. The sharp drop in sensitivity of the \( u\)-band filter at longer wavelengths ensures that objects at \( z \sim 3.42 \) (where transmission of the filter is less than 0.5 per cent) are free from flux contamination by non-ionizing flux longward of 912 Å (i.e. the Ly\( \alpha \) forest part). The shape of the transmission curve of the \( u \) filter is derived taking into account transmission from the CFHT primary mirror (black dotted line), MegaPrime optics (dotted–dashed line), and CCD quantum efficiency (dashed line). The blue dotted line represents median atmospheric extinction in mag/airmass units at the Mauna Kea site (Buton et al. 2012).

The MegaCam (Boulade et al. 2003) wide-field optical imager mounted on the CFHT (Canada–France–Hawaii Telescope) telescope is used during CLAUDS observations. The quantum efficiency of MegaCam ranges from 35 to 60 per cent at 3500 – 4000 Å, which makes MegaCam the most sensitive wide-field imager in this blue-optical wavelength range on any current telescope. The data from the new CLAUDS \( u\)-band filter used in this work in comparison with older \( u^\prime \) filter have two important advantages for our research; the new \( u \) filter probes bluer wavelengths and has a sharp cut-off at \( \sim 4000 \) Å and no red leak, whereas the older \( u^\prime \) filter probes redder wavelengths and has red leak at \( \sim 5000 \) Å.

Our aim is to test the hypothesis that LCGs at \( z \geq 3.5 \) with measurable LyC emission, have colours that reside in locations outside of the LBG colour-selection region as proposed by Cooke et al. (2014), see Section 2.2 for more details. The lowest secure redshift at which we can select our candidates is \( z = 3.42 \) as defined by the shape of the Megacam new \( u\)-band filter transmission curves. Above this redshift, only the LyC part of the spectrum is probed by the \( u \) band, with no chance of contamination by the Ly\( \alpha \) forest part (912 Å < \( \lambda \) < 1216 Å), even if the spectroscopic redshift is overestimated by \( \Delta z \sim 0.1 \). Fig. 1 shows the transmission curves of the \( u\)-band filter overlaid on an LBG spectrum in rest-frame wavelengths, alongside the transmission curves for the \( \text{Subaru} \) HSC \( r \) and \( i \) filters. In addition, the transmission curves of the CFHT primary mirror, MegaPrime optics, CCD quantum efficiency, and the atmospheric extinction for the Mauna Kea site are also shown.

The \( \text{Subaru} \) HSC is a wide-field camera mounted on the \( \text{Subaru} \) 8.2-m telescope (Hawaii, Mauna Kea site). The goal of the HSC survey is to observe high latitude fields like COSMOS in multiple photometric broad-band (\( g, r, i, z, \) and \( y \)) and narrow-band filters (Aihara et al. 2018b). The filter depths based on PDR1 are 27.4,
LyC flux. Thus, expected to reside off this plot using the standard LBG selection criteria.

5, redshift, respectively. where triangles, squares, circles and stars correspond to are shown by solid line curves. The expected redshifts are marked in black, 27.3, 27, 26.4, and 25.6 mag for g, r, i, z, and y respectively, (5σ depth for a point source; Aihara et al. 2018b). Table 1 summarizes the photometric filter properties used in this work.

### Table 1. Photometric filter characteristics.

| Filter | λ_eff (Å) | FWHM (Å) | 5σ depth (mag) | Survey       |
|--------|-----------|-----------|----------------|--------------|
| u      | 3538      | 860       | 27.2           | CLAUDS \(^a\) |
| g      | 4754      | 1395      | 27.4           | Subaru HSC \(^b\) |
| r      | 6175      | 1503      | 27.3           | Subaru HSC \(^b\) |
| i      | 7711      | 1574      | 27.0           | Subaru HSC \(^b\) |
| z      | 8898      | 766       | 26.4           | Subaru HSC \(^b\) |
| y      | 9762      | 783       | 25.6           | Subaru HSC \(^b\) |

\(^a\) Limiting magnitude estimated with 2 arcsec diameter aperture.

\(^b\) Subaru HSC limiting magnitudes in PDR1 are for point sources.

#### Figure 2. Colour–colour u - g versus g - i plot, the LBG selection region for z ~ 2.7–3.4 galaxies is shown in grey. LBG galaxies at z > 3.4 are expected to reside off this plot using the standard LBG selection criteria. The presence of z > 3.4 galaxies on this plot suggests the presence of LyC flux. Thus, u - g versus g - i colour–colour plane provides a good means to identify LyC leaking galaxies. Dotted blue, green, yellow, and red lines are evolutionary tracks of the four composite spectra adopted from Shapley et al. (2003) with different EW(Ly\(\alpha\)) and without LyC flux. Evolutionary tracks after adding LyC flux to composite spectra in different ratios (R_{LyC} = 1, 2, 5, 10, and 20 per cent, R_{LyC} is defined in Section 4) are shown by solid line curves. The expected redshifts are marked in black, where triangles, squares, circles and stars correspond to z ~ 3.5, 4, 4.5, and 5, redshift, respectively.

27.3, 27, 26.4, and 25.6 mag for g, r, i, z, and y respectively, (5σ depth for a point source; Aihara et al. 2018b). Table 1 summarizes the photometric filter properties used in this work.

#### 2.2 Parent sample

To date, high-redshift galaxies with escaping LyC photons have usually been selected using the Lyman-break technique. It is recognized that during LBG selection a significant number of the galaxies that are not classified as LBGs (outside the standard LBG selection region) and that are at z ≥ 3.5. This approach allows us to test predictions from Cooke et al. (2014).

In this work, we utilize spectroscopic redshifts from several surveys in the literature that cover the COSMOS field. By selecting objects with redshifts z ≥ 3.5 that have high-quality flags, determined by different surveys, we ensure that any flux detection in the CLAUDS u-band imaging is a clean LyC detection. This means that the flux observed for the object is not contaminated by the non-ionizing radiation.

To create a parent sample that will be examined in Section 2.3 for LyC radiation, we first select all objects with spectroscopic redshifts z ≥ 3.5 from the D10K, VUDS, and 3D-HST catalogues. The selected objects have high-quality spectrum flags that correspond to ≥ 75 per cent probability that the reported redshift is correct.

The initial selection results in 407 objects and this parent sample is presented in Table 2 and the redshift distribution is shown in Fig. 3. Most of the redshifts in our parent sample, 361 objects (contributing 89 per cent to the total sample used here) with z ≥ 3.5, are from the D10K survey. The D10K survey selects objects from a variety of input catalogues based on multiwavelength observations and, importantly, have different selection criteria. The full D10K survey uses multislit spectroscopy that covers the wavelength range ~5500–9800 Å and objects are identified up to z ~ 6. More details on the observations, target selection, and reduction can be found in Hasinger et al. (2018). From the VUDS spectroscopic survey we were able to extract 40 objects (10 per cent of the total sample) that have z ≥ 3.5. VUDS spectroscopy covers the 3650–9350 Å wavelength range and targets objects at all redshifts to z ~ 6. More details on the observations, target selection, and reduction can be found in Le Fèvre et al. (2015). From the available grism data that are a product of the 3D-HST survey, we extract six objects.

### Table 2. Parent sample of galaxies with spectroscopic redshifts.

| Survey | Number of objects | Percentage of parent sample |
|--------|-------------------|----------------------------|
| D10k   | 361               | 89                         |
| VUDS   | 40                | 10                         |
| 3D-HST | 6                 | 1                          |
| Total  | 407               | 100                        |

#### Figure 3. Redshift distribution of the parent sample, where the width of the bins is Δz = 0.2. The peak at z ~ 5.7 is due to narrow-band LAE-selected sample.
(1 per cent) with \( z \geq 3.5 \). For more details about 3D-HST grism spectroscopy, we refer the reader to Momcheva et al. (2016).

Next, the parent sample is cross-matched with the Subaru HSC catalogue to obtain photometry of the sources within 0.5 arcsec using the TOPCAT software (Taylor 2005), resulting on 375 object matches (32 objects not matched). The aperture centres for the 375 objects are defined by adopting coordinates from the Subaru HSC catalogue. For the other 32 candidates, we use the coordinates from the spectroscopic catalogues.

### 2.3 Subsamples

With the parent sample now defined, the next step is to perform aperture photometry. We perform 1.2 arcsec diameter circular aperture photometry with the PYTHON ASTROPY package PHOTUTILS designed to detect and perform photometry of astronomical sources. Photometry on the CLAUDS \( u \)-band images resulted in positive flux detection in 22 candidates with \( S/N > 3 \) and 151 candidates with \( S/N < 3 \). For 234 candidates the estimated LyC flux was negative. For this work, we will continue our analysis on the objects with reported positive flux for a total of 173 candidates (22 with \( S/N > 3 \) and 151 with \( 0 < S/N < 3 \) ) and from 234 candidates with negative flux we will use only candidates that show no other object inside a radius of 1 arcsec from the candidate.

In the next phase of the sample selection, we visually inspect all 173 candidates with positive flux. Thumbnails of size 15 arcsec \( \times \) 15 arcsec are created in all available filters and all candidates are checked for LyC signal using CLAUDS \( u \)-band image with \( \sim 0.2 \) arcsec pixel \(^{-1} \) resolution and median seeing of 0.80–0.85 arcsec (Sawicki et al. 2019). The Subaru HSC images in the \( g, r, i, \) and \( z \) filters are also inspected for flux contamination from nearby bright objects and star spikes. Previous research by Vanzella et al. (2010a) and Siana et al. (2015) indicates that the probability of contamination by foreground objects should not be neglected and depends on seeing conditions, aperture size, limiting magnitude, and increasing redshift. Based on the Vanzella et al. (2010a) results, we would expect that the fraction of contaminated candidates in our sample would be \( \sim 3 – 15 \) per cent in the case of the 0.5–1 arcsec seeing, respectively. To decrease the fraction of candidates contaminated by foreground objects, we use high-resolution HST imaging in the F125W and F160W (Koekemoer et al. 2011; Grogin et al. 2011; Skelton et al. 2014), F140W (Skelton et al. 2014), F336W and F435W (Pritchard et al. in preparation), F606W (Koekemoer et al. 2011; Grogin et al., 2011), and F814W bands. The HST F814W and F606W filter images have \( \sim 0.03 \) arcsec pixel resolution, while images in other HST filters have \( \sim 0.06 \) arcsec pixel resolution. Imaging in all mentioned HST filters is not available for every object due to different areas of HST imaging coverage, however, F814W is available for 95.6 per cent (389) of our sample of 407 objects. The same strategy of the visual inspection is adopted for the 234 candidates with negative flux to select only candidates without any source of radiation within 1 arcsec radius. We classify 53 candidates with negative flux as not contaminated by flux from the nearby galaxies or other sources of radiation, with 181 candidates are not suitable for further analysis due to external flux contamination.

For the purpose of the visual classification, the five subsamples are defined as follows:

(i) **Detection**: these are the objects that are considered clean, non-contaminated LyC emitters. They are defined as galaxies with no evidence of another object within a 1 arcsec radius and the signal of the candidate in any band (\( u, g, r, i, z \), or any available \( HST \) band) is not contaminated by flux leakage from nearby bright objects or star spikes and there is no evidence of an intervening object in the spectra. We identify two objects (0.5 per cent of our parent sample) in this subsample with \( S/N \geq 3 \).

(ii) **Detection close pairs**: these are objects where the candidate shows LyC flux but another object is located within a 1 arcsec radius and does not appear to contaminate the LyC radiation from candidate. We identify five objects (1.2 per cent of our parent sample) in this subsample with \( S/N \geq 3 \).

(iii) **Non-detection**: these are non-contaminated objects (i.e. no object is detected within a 1 arcsec radius) for which LyC radiation is not detected. We identify 35 objects with positive measured flux and 53 objects with negative measured flux, in total 87 objects (21.6 per cent of our parent sample) in this subsample.

(iv) **Multiple objects**: these are the objects that appear to be more than one galaxy within a 1 arcsec radius and without LyC flux from the candidate, and are either mergers or are contaminated by low-redshift interloper(s). It is difficult to distinguish the two types from images alone. By default we include any object without HST coverage in this group, as we are not able to determine whether the source is contaminated by another object within the 1 arcsec radius. Objects from this subsample will not be considered further in this work. We identify 118 objects (29 per cent of our parent sample) in this subsample.

(v) **Flux contaminated**: these are the objects that are contaminated by flux from nearby bright galaxies, stars or star spikes in any of the examined filters. Objects in this subsample will not be used for any kind of analysis. We identify 13 objects (3.2 per cent of our parent sample) in this subsample.

(vi) **Negative flux − contaminated**: these are the objects that appear to be more than one galaxy within a 1 arcsec radius and measured flux from the candidate is negative. We identify 182 of these objects (44.7 per cent of our parent sample)

The final outcome from visual classification of the parent sample is summarized in Table 3 where all candidates in the Detection and Detection close pairs subsample have \( S/N > 3 \) and all non detection have \( S/N < 3 \).

### 3 SPECTRAL CONFIRMATION

With the methodology now described we will proceed with a description of carefully inspecting the subsamples, with the goal of creating a non-contaminated sample that is as clean as possible from low-redshift objects. We inspect the 1D and 2D spectra (if available) to double check the reported spectroscopic redshifts and carefully examine the spectrum for any possible foreground contamination. The slit position (if available) of the candidate is also compared with the position of the dispersion line in the 2D spectrum to confirm that

| Subsample | Number of objects | Percentage |
|-----------|-------------------|------------|
| Detection | 2 | 0.5 |
| Detection close pairs | 5 | 1.2 |
| Non-detection | 87 | 21.4 |
| Multiple objects | 118 | 29 |
| Flux contaminated | 13 | 3.2 |
| Negative flux − contaminated | 182 | 44.7 |
| Total | 407 | 100 |
the reported redshift belongs to the candidate object. Furthermore, to help rule out the possibility that our candidates are AGNs/QSOs, the spectra are examined for AGN far-UV emission-line signatures and Detection and Detection close pairs are cross-matched with the publicly available XMM–Newton X-ray (Cappelluti et al. 2009) and Chandra COSMOS legacy (Marchesi et al. 2016) catalogues within 0.5 arcsec. The results of the spectral confirmation are summarized in Table 4 and described in the following Sections 3.1 and 3.2, where two quality groups are created. For objects classified as quality 1 (q1), we were able to confirm their reported redshifts, and for quality 2 (q2), we were unable to solidly confirm reported redshifts.

### 3.1 Detection subsample

After visual selection and spectroscopic examination, our Detection subsample contains two candidates that show possible non-contaminated LyC flux. For both candidates 1D and 2D spectra are re-inspected to confirm redshifts and rule out possible contamination by low-redshift interlopers. The EW for all candidates are measured in the same way as described in Cassata et al. (2015) using the IRAF (Tody 1986) SPLOT tool from NOAO.ONEDSPEC package.

By cross-matching coordinates with the XMM–Newton X-ray and Chandra legacy catalogue, we did not characterize either candidate as an AGN. Individual analysis of the available 1D and 2D spectra from the literature, as well as slit positions and multiband imaging data revealed the following:

(i) Candidate ID 394: the 1D and 2D spectra from the VUDS survey were available for analysis. Analysing 1D and 2D spectra, we were unable to claim any spectroscopic feature that would indicate reported redshift, so we exclude this candidate from further analysis.

(ii) Candidate ID 421: the spectrum from the VUDS survey was available for analysis. From the spectrum we were able to identify several spectroscopic features consistent with the reported redshift from the VUDS survey, such as a Lyα forest break, Lyβ absorption, and potential ISM absorption features, but conservatively cannot claim a solid redshift confirmation. The spectrum is presented in Appendix B.

In Section 5, we will refer to the object ID 421 as a second quality (q2) candidate. Thumbnails of the Detection subsample in the CLAUDS u and HST F814W bands are presented in Fig. 4.

### 3.2 Detection close pairs subsample

For the Detection close pairs subsample, the same analysis procedure is applied as for the Detection subsample. Thumbnails of the Detection close pairs subsample in the CLAUDS-u and HST F814W bands are presented in Fig. 5. We visually classify five objects as

![Figure 4. Thumbnails 15 arcsec × 15 arcsec in size for the Detection subsample. Candidates are shown in two bands: CLAUDS-u and HST F814W. The red circle is 2 arcsec diameter aperture. Next to the ID number of the candidate in brackets are quality group of the objects (q1 or q2) or x which means that candidate is rejected.

Detection close pairs that potentially show LyC flux. Individual analysis of the available 1D and 2D spectra from the literature, as well as slit positions and multiband imaging data, revealed the following information on the Detection close pairs candidates:

(i) Candidate ID 1 has features of an Lyα emitter, namely a relatively strong, asymmetric emission feature in the 1D and 2D spectrum typical of Lyα lines at high redshift and a higher flux level in the UV continuum redward of Lyα as compared to blueward (i.e. the Lyα forest). This object will be noted as first quality (q1) candidate.

(ii) Candidates ID 326 and 330 have somewhat asymmetric emission lines, consistent with those observed for LBGs, and are probably Lyα (see Appendix B). The continua are faint or not detected in the spectra. For these candidates we were not able to securely confirm their reported redshifts and in Section 5 these objects will be noted as second quality (q2) candidates.

(iii) Candidate ID 368 has detected Lyα with rest-frame EW(Lyα) ~ 25 Å. The Lyα spectral profile, appropriately strong Lyα forest break, UV continuum profile, and ISM line absorption features are consistent with a z = 3.64 SFG. Further inspection of the spectra did not reveal any other lines that could be related...
to low-redshift objects. Candidate ID 368 was part of the clean sample in Marchi et al. (2017) where they estimate the LyC signal from galaxies at $z \sim 4$ using VIMOS spectroscopy and available HST imaging. Furthermore, Marchi et al. (2017) do not report any single solid detection of individual candidates. On the contrary, in this work using CLAUDS $u$-band photometry we detect an LyC signal from candidate ID 368, and its $z = 3.64$ has now been additionally confirmed with MOSFIRE spectroscopy where the nebular emission lines [OIII] doublet and H$\beta$ are detected (Bassett et al., in preparation). Beside nebular lines, the MOSFIRE spectrum reveals no other lines, which additionally confirms that detected LyC signal is not contaminated by low-redshift interlopers. VUDS and MOSFIRE spectra are shown in Appendix B. In Section 5, we will refer to the candidate id 368 as a first quality candidate (q1).

(iv) Candidate ID 359 has two or more independent emission lines indicating a merger system or contamination by a low-redshift object and both objects fall in the slit. We label the bottom object ‘object A’ and upper object ‘object B’ in Fig. 5. In the 2D spectrum, the two emission lines are visible, as well as two faint continua which is in agreement with the positions of the objects in the slit. From the 2D spectrum, we conclude that the line at $\sim 7747$ Å is related to object A, whereas the line at $\sim 8167$ Å is related to object B. If both detected lines are Ly$\alpha$ the estimated redshifts would be $z \sim 5.37$ and $\sim 5.72$ for objects A and B, respectively. Since the spectrum has a low S/N (signal-to-noise ratio) and the two objects are close to each other it is difficult to extract them separately and to check for the expected drop in flux blueward of the Ly$\alpha$ line. As a result, we choose to remove this candidate from the further analysis.

To conclude this section on spectral confirmation, we provide a flow chart in Fig. 6 where the complete selection process of our LyC candidates is summarized. In total, we find two q1 and three q2 candidates.

4 CALCULATION AND EXPRESSION OF THE LYC ESCAPE FRACTION

One of the least constrained parameters in both observational and theoretical studies of reionization is the escape fraction of LyC photons ($f_{\text{esc}}$). The escaping LyC photons produced by O and B stars in SFGs are referred to as the absolute $f_{\text{esc,abs}}$ or relative $f_{\text{esc}}$ or $f_{\text{rel,esc}}$. The quantity $f_{\text{rel,esc}}$ is defined as the fraction of the ionizing photons that escape without being absorbed by interstellar medium (ISM) or circumgalactic medium (CGM) into the IGM. Direct measurements of $f_{\text{esc,abs}}$ are not possible since the intrinsic ratio of the ionizing and non-ionizing UV photons is difficult to determine. Moreover, the measurement of the ionizing and non-ionizing UV radiation is severely affected and suppressed by the high IGM opacity towards higher redshifts and dust attenuation. We therefore have to model these missing parameters. Modelling SEDs of distant SFGs by using different stellar population synthesis models and taking into account dust extinction can produce insights in the intrinsic properties of the galaxies at high redshifts. The quantity estimated to help observations of LyC emitters is $f_{\text{rel,esc}}$ (e.g. Steidel et al. 2001), and is defined as:

$$f_{\text{rel,esc}} = \frac{(F_{\text{LyC}}/F_{1500})_{\text{obs}}}{(L_{\text{LyC}}/L_{1500})_{\text{int}}} \exp(\tau_{\text{LyC,IGM}}),$$

where $(F_{\text{LyC}}/F_{1500})_{\text{obs}}$ is the observed rest-frame LyC to UV flux density, $(L_{\text{LyC}}/L_{1500})_{\text{int}}$ is the intrinsic ratio of the galactic...
ionizing (LyC) to non-ionizing (UV) luminosity density, and $\tau_{\text{LyC}}$ is the redshift-dependent attenuation of LyC photons due to intergalactic neutral hydrogen along the line of sight. From models, we assume $(L_{\text{LyC}}/L_{1500})_{\text{int}}$ varies in the range of $\sim 0.1$–0.9 and depends on several galactic parameters like star formation history, stellar initial mass function (IMF), stellar age, and metallicity. The attenuation factor $\tau_{\text{LyC}}$ is usually determined by analytic models or MC simulations (e.g. Inoue et al. 2014; Steidel et al. 2018). Finally, we directly measure $(F_{\text{LyC}}/F_{1500})_{\text{obs}}$ in this work.

We use CLAUDS $u$ band to measure $(F_{\text{LyC}})_{\text{obs}}$, corresponding exclusively to the LyC flux over the rest-frame wavelength range. To probe $(F_{1500})_{\text{obs}}$, rest-frame non-ionizing UV flux ($\sim 1500$ Å), we use Subaru HSC photometry PDR1 in $i$, $z$, and $y$ bands, depending on the redshift of the galaxy. For the objects in the range $3.5 < z < 4.5$, we use $i$ band, $z$ band is used at the range $4.5 < z < 5.5$, and $y$ band in the $5.5 < z < 6.5$ range. The non-ionizing UV flux is measured in the same aperture size (1.2 arcsec diameter) and centred on the same coordinate as for LyC flux from CLAUDS $u$-band images. In this way, we ensure that the same parts of the galaxy are probed. The results from equation (1) can be directly converted to $f_{\text{esc}}^{\text{abs}}$ by following the equation proposed by Inoue et al. (2005) and Siana et al. (2007):

$$f_{\text{esc}}^{\text{abs}} = f_{\text{esc}}^{\text{rel}} \times 10^{-0.4k_{\lambda}E(B-V)},$$

(2)

where $k_{\lambda}$ is the reddening law that describes how the chosen dust model affects the UV radiation at particular wavelengths (here we will use $k_{\lambda} = 10.33$ for a Calzetti reddening law; Calzetti 1997) and $E(B - V)$ is the total dust attenuation or reddening. For the purpose of our work, we use $E(B - V)$ values from the publicly available catalogue of Laigle et al. (2016). By correcting $f_{\text{esc}}^{\text{rel}}$ values for internal dust attenuation we are able to obtain rough estimates for $f_{\text{esc}}^{\text{abs}}$.

The only measured value in equation (1) is $(F_{\text{LyC}}/F_{1500})_{\text{obs}}$. Another way to express this quantity is the so-called relative observed fraction or $R_{\text{obs}}(\lambda)$ proposed by Cooke et al. (2014):

$$R_{\text{obs}}(\lambda) = \frac{F_{\text{LyC}}}{F_{\text{UV}}},$$

(3)

where $F_{\text{LyC}}$ is the observed ionizing radiation flux integrated over the filter probing the LyC (here, the CLAUDS $u$ band) and $F_{\text{UV}}$ is the observed non-ionizing UV radiation near 1500 Å (here, the Subaru, $i$ or $z$ bands). In the literature $R_{\text{obs}}(\lambda)$ is usually referred to as the flux density ratio, $(f_{\text{1500}}/f_{\text{900}})_{\text{abs}}$. The advantages of using $R_{\text{obs}}(\lambda)$ in applications like this is that it is not model dependent, derived directly from the observations. In addition, it results in an arguably more practical value (i.e. per cent LyC flux) and the typically smaller errors on the UV continuum are in the denominator. The value $R_{\text{abs}}(\lambda)$ is advantageous here, as we are using only data from filter observations and in comparison with other galaxy quantities that are also derived from the observations [e.g. magnitude, colour, and EW(Ly$\alpha$)].
Figure 7. The age of the stellar population plotted against estimated $f_{\text{esc}}$ for our five LCG candidates, where age of the stellar population also can be interpreted as different $(L_{\text{LyC}}/L_{1500})_{\text{int}}$. Here, we adopt the mean $\tau_{\text{LyC}}^{\text{IGM}}$ from Inoue et al. (2014, see the text for more details). The red vertical line marks where $f_{\text{esc}} = 100$ per cent and horizontal line shows at which point O stars end their production of LyC photons. Blue circles, pink squares, and green triangles represent the default imf135 model from BPASSv2.2 with subsolar, solar, and supersolar metallicities, respectively. For the candidate ID 330, the estimated $f_{\text{esc}}$ is beyond 100 per cent (the points are outside the graph).

4.1 Attenuation of the IGM – $\tau_{\text{IGM}}^{\text{LyC}}$

For the purpose of this work, the mean $\tau_{\text{IGM}}^{\text{LyC}}$ is estimated by using the results from the updated analytic model for attenuation presented by Inoue et al. (2014). Here, $\tau_{\text{IGM}}^{\text{LyC}}$ is derived as the weighted average across the whole $u$-band filter. In that way, we are taking into the account the transmission variation of the filter. We find that adopting the mean $\tau_{\text{IGM}}^{\text{LyC}}$ is not always the best strategy for individual objects. In cases where we have a confirmed detection of LyC flux, the probability that this line of sight has a higher IGM transparency than the mean is not negligible. In these situations, using the mean $\tau_{\text{IGM}}^{\text{LyC}}$ can result in an overestimation of $f_{\text{esc}}$. For example, the probability of a clean line of sight, where optical depth is less than unity and free from Lyman limit systems (LLSs), at 900 Å (source rest frame) is estimated to be $\sim 70$ per cent for objects at $z = 3$, and $\sim 20$ per cent for $z = 4$ (Inoue & Iwata 2008).

To account for the fact that any detection of LyC radiation at $z > 3.6$ is likely to arise from a line of sight that is more transparent than the average transmission, we also adopt the maximum transmission ($1 - D_0$) estimated for the redshift range $2.7 \leq z \leq 5$, from Steidel et al. (2018). The quantity $D_0$ is defined as a mean depression in the rest-frame continuum interval $920$–$1015$ Å caused by Lyman line blanketing (Oke & Korycansky 1982). As detailed in their work, $(1 - D_0)$ is a close approximation for the maximum IGM+CGM transparency expected at given redshift that is estimated in rest-frame wavelength interval $880$ Å $\leq \lambda_{\text{rest}} \leq 900$ Å. The caveat to this approach is, for our objects, that we are probing the LyC region shortward of $880$ Å. Thus we are assuming that ionizing radiation shortward of $912$ Å is more or less equally affected by IGM+CGM. In this case adopting $(1 - D_0)$ as a correction factor for IGM attenuation will give us lowest possible predictions for $f_{\text{esc}}$.

In this work, we are dealing with single lines of sight and estimating the actual value of $\tau_{\text{IGM}}^{\text{LyC}}$ for each case is extremely complicated and beyond the scope of this paper. This topic will be discussed in a forthcoming paper by Bassett et al. (in preparation). Here, we adopt these two stochastic extremes for IGM transmission, the mean $\tau_{\text{IGM}}^{\text{LyC}}$ and $(1 - D_0)$, to bracket the upper and lower limits for $f_{\text{esc}}$, respectively, with each limit being an unlikely case.

4.2 The intrinsic luminosity ratio $(L_{\text{LyC}}/L_{1500})_{\text{int}}$

The intrinsic luminosity ratio is the most poorly constrained quantity in equation (1). Since we are unable to put solid constraints on the intrinsic luminosity of the galaxy from direct observations, in most cases in the literature the $(L_{\text{LyC}}/L_{1500})_{\text{int}}$ ratio is estimated from using stellar population synthesis models. In this work,
Figure 8. Same as Fig. 7, instead of the mean $\tau_{\text{LyC}}$ values, we adopt a close approximation to the maximum transmission, $(1 - D_b)$ from Steidel et al. (2018) in the interval $880 \, \text{Å} \leq \lambda_{\text{rest}} \leq 900 \, \text{Å}$.

Figure 9. $u - g$ colour as a function of $u_{\text{mag}}$. Five LyC Candidates with q1 (blue squares) and q2 (green circles) are plotted. Colour coded hex bins present distribution of objects from the CLAUDS catalogue.

we adopt results from Binary Population and Spectral Synthesis models (BPASSv2.2, Stanway & Eldridge 2018). To estimate $(L_{\text{LyC}}/L_{1500})_{\text{int}}$, we use the default $\text{imf135}_300$ binary population single instantaneous burst model over the mass range $0.1 - 300 \, M_\odot$.\(^2\) With three different metallicities subsolar, solar, and supersolar (zem5, z020, and z040) for more details see Stanway & Eldridge (2018) and the BPASS manual.\(^3\) We calculate the $(L_{\text{LyC}}/L_{1500})_{\text{int}}$ quantity for every single candidate separately, as it depends on the time since the onset of star formation. For this purpose, we are using three different types of synthetic spectra with subsolar (zem5), solar (z020), and z040 metallicities in the $10^6 - 10^9$ age span. We first convert the flux of the BPASS model synthetic spectra from solar luminosity per angstrom to luminosity densities (erg $s^{-1}Hz^{-1}$). We then normalize the synthetic extreme UV spectral region (200–1750 Å) by the mean flux value in the range 1450–1525 Å. Finally, the value of $(L_{\text{LyC}}/L_{1500})_{\text{int}}$ is estimated based on the coverage of the $u$-band filter for every candidate. For example, for an LyC candidate at $z \sim 3.6$, we take the average value in the 668–868 Å range (rest frame of the candidate) corresponding to the FWHM of the $u$-band filter for that redshift.

5 RESULTS AND DISCUSSION

From here on, we present our LyC candidates in two quality groups q1 (2 objects) and q2 (3 objects) for a total of five objects. The quality group assignment is based on the quality of the candidates in the images, their spectroscopic redshift, and lack of low-redshift contamination evidence (Section 3). Table 5 presents the observed $R_{\text{obs}}$, the rest-frame LyC wavelength coverage by the CLAUDS catalogue.

\(^2\)The adopted IMF is based on the Kroupa, Tout & Gilmore (1993).

\(^3\)https://drive.google.com/file/d/1ImqPuFTYLQ7k/view

Downloaded from https://academic.oup.com/mnras/article-abstract/494/4/4986/5818351 by California Institute of Technology user on 25 June 2020
5.1 Estimates of \( f_{\text{esc}} \)

From the definition of \( f_{\text{esc}} \), introduced in Section 4, it is clear that \( f_{\text{esc}} \) estimates strongly depend on the accuracy of the modelled parameters. Under these circumstances, it is difficult to constrain the escaping flux from the candidates when there are two ‘free parameters’ that span 0–1 for \( \tau_{\text{IGM}} \) and 0.1–0.9 for \( (L_{\text{LyC}}/L_{1500})_{\text{int}} \). Because of this difficulty, we decide to estimate \( f_{\text{esc}}^{\text{abs}} \) by using different \( \tau_{\text{IGM}} \) and \( (L_{\text{LyC}}/L_{1500})_{\text{int}} \) quantities. As a result, \( f_{\text{esc}}^{\text{abs}} \) values in this work are presented as a range of values, rather than as a single value.

5.1.1 Estimates of \( f_{\text{esc}} \) based on the mean \( \tau_{\text{IGM}} \)

As discussed, \( \tau_{\text{IGM}} \) depends on the redshift and distribution of matter along the line of sight. The \( \tau_{\text{IGM}} \) is estimated as a weighted average across the entire \( u \)-band wavelength range. The IGM attenuation is estimated by using results from the analytical methods from Inoue et al. (2014), where the IGM transmission is estimated by averaging over 100,000 lines of sight. Some of the sightlines may intersect relatively rare LLSs and, in that case, LyC photons are severely attenuated. An example of how LyC photons are affected when an LLS is in the line of sight and when one is absent is shown in fig. 6 of Inoue & Iwata (2008).

It is important to note that the redshifts of the five potential LCGs \((q_1 \text{ and } q_2)\) reported in this work span from \( z \sim 3.5 \) to \( z \sim 5 \), close to the period when reionization ends. Thus, in most cases we are probing much bluer parts of the LyC (see Table 5) than the 880 Å < \( \lambda_{\text{rest}} \) < 912 Å range most often studied in the literature.

We estimate \( f_{\text{esc}}^{\text{abs}} \) for our candidates after adopting the mean \( \tau_{\text{IGM}} \) across the entire \( u \)-band using analytical models and \( (L_{\text{LyC}}/L_{1500})_{\text{int}} \) derived from BPASS. The results are presented in Fig. 7 where the age of the stellar population, that spans from 10\(^6\) to 10\(^9\) yr, is plotted against estimates on \( f_{\text{esc}}^{\text{abs}} \). Blue circles, pink squares, and green triangles represent the same default IMF135_300 model from BPASSv2.2 with different metallicities subsolar, solar, and supersolar, respectively. The vertical red line on the graph marks where \( f_{\text{esc}} = 100 \) per cent. LyC radiation is mostly produced by short-lived massive O-type stars, whose lifetimes are \( \sim 10^7 \) yr. The horizontal red line marks the period when we expect the production of the LyC photons from O-type stars is terminated. As a result, the most relevant part of the plot is likely the 10\(^6\) to 10\(^9\) yr time range. We estimate \( f_{\text{esc}}^{\text{abs}} \) at different young stellar population ages, which can also be interpreted as different \( (L_{\text{LyC}}/L_{1500})_{\text{int}} \). As the stellar population ages, the ratio between ionizing LyC and non-ionizing UV photons decreases.

For the candidate ID 330, \( f_{\text{esc}}^{\text{abs}} \) > > 100 per cent. This result can be an indication that the LyC flux from our candidate is contaminated by a low-redshift interloper or that the estimated redshift is lower than 3.4. Another interpretation, since candidate ID 330 is classified as detection-close pairs, is that this is a merger system. Mergers are not taken into account by BPASS population synthesis models and in these scenarios, LyC photons could be produced by fast radiative shocks, the interaction of the clumps, or accretion processes (Dopita et al. 2011; Wyihte, Mould & Loeb 2011).

From Fig. 7, all of the candidates show \( f_{\text{esc}}^{\text{abs}} \) > > 10 per cent after the O stars start to produce LyC photons. If we assume that, at the beginning of their evolutionary paths, O stars are embedded in clouds of dust and gas, it is more likely to expect that \( f_{\text{esc}}^{\text{abs}} \) starts from < 0.15 and gradually increases. As evolution progresses, more ionizing UV radiation is emitted and the material around the stars gets pushed away and ionized. This can lead to the formation of the ‘clean’ paths or holes around systems that produce LyC photons and through these directions they can freely escape into the IGM (Zackrisson, Inoue & Jensen 2013).

5.1.2 Estimates of \( f_{\text{esc}} \) based on higher IGM transparency

Our other approach to estimate \( f_{\text{esc}}^{\text{abs}} \) is based on the assumption that the detected \( u \)-band flux in these high-redshift sources indicates that our search, with its current sensitivity, is most likely biased toward lines of sight with low H\(_I\) densities or free from LLSs. Because of this, we consider whether the assumption to use the mean IGM transmission for those objects with directly detected

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Table 5. Observed and modelled properties for the five LCGs \( q_1 \) and \( q_2 \) candidates.

| ID | \( z_{\text{spec}} \) | \( \tau_{\text{IGM}} \) | \( (1 - D_h)^{b} \) | \( \text{LyC} \lambda_{\text{rest}}(\AA)^{c} \) | \( R_{\text{obs}}(\lambda)^{d} \) | \( E(B - V)^d \) | \( f_{\text{esc}}^{\text{abs}} \) (per cent) |
|----|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|
| 1  | 4.28            | 0.006           | 0.334           | 568–763         | 0.11 ± 0.03     | 0.4             | \( \gtrsim 5 - 73 \) |
| 326| 3.57            | 0.097           | 0.565           | 657–882         | 0.16 ± 0.04     | 0.3             | \( \gtrsim 4 - 15 \) |
| 330| 5.09            | 0.00016         | 0.233           | 493–662         | 0.15 ± 0.04     | 0               | > 100           |
| 368| 3.64            | 0.08            | 0.565           | 647–869         | 0.12 ± 0.04     | 0.1             | \( \gtrsim 30 - 93 \) |
| 421| 3.60            | 0.09            | 0.565           | 652–876         | 0.17 ± 0.04     | 0.2             | \( \gtrsim 8 - 47 \) |

\(^{a}\) Mean IGM transmission estimated from Inoue et al. (2014).

\(^{b}\) Close approximation to the maximum IGM + CGM transmission (Steidel et al. 2018).

\(^{c}\) Rest-frame LyC probed by CLAUDS \( u \) band.

\(^{d}\) Adopted from Laigle et al. (2016).

\(^{e}\) Possible low-redshift galaxy or contaminated by low-\( z \) interloper.
LyC radiation is appropriate. The influence of the IGM is also discussed by Vanzella et al. (2010b) and their findings indicate that in some cases, transmission along the lines of sight drops to zero blueward of the redshift of the LLS, but there are cases where signal from the source is transmitted down to ~700 Å in agreement with our observations. If the IGM transmission of the line of sight is assumed to be higher than the mean value at a particular redshift, \( f_{\text{esc}} \) of the source can be smaller than the values shown in Fig. 7.

We have shown in Section 5.1.1 that the LyC escape fraction may be overestimated if we use the mean IGM transmission to correct for IGM attenuation. Therefore, we use IGM transmission values that describe lines of sight with higher transparency for LyC radiation. For this purpose, we adopt the estimated maximum IGM+CGM transmission values from Steidel et al. (2018). A shortcoming of this approach is that the estimated IGM+CGM transmission is only available for the rest frame 880 Å \( \leq \lambda_{\text{rest}} \leq 900 \) Å interval and the chance for not having high column density absorber (\( \tau \)) at the redshifts of our candidates is less likely. Therefore, the estimated range of \( f_{\text{esc}} \) for our candidates can be interpreted only as lower limits. Also for our redshift range, we only have IGM+CGM transmission estimates for discrete redshift bins \( z = 3.5, 4, 4.5, \) and 5. The approximations of the maximum IGM transmission \( ((1-D_b)) \) that we adopt for each candidate in our sample are summarized in Table 5. The \( f_{\text{esc}} \) results based on the \( (1-D_b)) \) values are shown in Fig. 8, where markers and axes are the same as Fig. 7.

In contrast to the previous estimates where mean IGM transmission is used, adopting the maximum transmission \( (1-D_b) \) leads to the lower \( f_{\text{esc}} \) values for both q1 and q2 candidates. The candidate ID 330 that belongs to the q2 sample with \( f_{\text{esc}} > 100 \) per cent can be explained as a low-redshift object or possible merger system as discussed in Section 5.1.1. Therefore, by taking into account only q1 objects, we can conclude that adopting the mean IGM transmission may not properly describe the extent to which LyC photons are attenuated in individual LyC emitters (this conclusion does not change if we also include q2 objects). On the other hand, the existence of very clean lines of sight is also less likely case at these given redshifts. Models, with different metallicities can produce variations in the estimated \( f_{\text{esc}} \) from a few percent in early evolutionary stages of the O stars up to ~10 per cent in late evolutionary stages, for both q1 and q2 candidates.

From the results presented in Figs 7 and 8 it is clear that even with a clean detection of the LyC flux it is extremely difficult to estimate the amount of LyC photons that are leaking into the IGM, as the two parameters, \( \tau_{\text{LyC}}^{\text{IGM}} \) and \( (L_{\text{LyC}}/L_{\text{LyC}4997})_{\text{int}} \) are almost impossible to measure directly and difficult to constrain with models or simulations for individual galaxies. By not knowing at least one of the model parameters more precisely our calculated \( f_{\text{esc}} \) values are rough estimates that range from a few percent up to ~90 per cent (candidate ID 330 excluded), in our cases, if we assume that O-type stars are the main producers of LyC photons. Therefore, as a final result we are adopting range in upper and lower limits estimated by using \( \tau_{\text{LyC}}^{\text{IGM}} \) and \( (1-D_b) \) respectively as a most plausible range for \( f_{\text{esc}} \) where adopted metallicity is zem5 (subsolar, Table 5).

5.2 Ly\( \alpha \) properties of the LCGs

In the last few years, the link between the properties of the Ly\( \alpha \) line, the escape of the Ly\( \alpha \) photons, and the escape of the LyC photons among low- and high-redshift galaxies has been explored by a number of independent studies. For example, the correlation between EW(Ly\( \alpha \)) and \( f_{\text{esc}} \) among low-redshift galaxies is reported by Verhamme et al. (2017), Steidel et al. (2018), and Fletcher et al. (2019) and an indication of a possible trend among EW(Ly\( \alpha \)) and \( R_{\text{abs}} \) is discussed by Marchi et al. (2017). In Fig. 10, we present rest-frame EW(Ly\( \alpha \)) as a function of estimated \( f_{\text{esc}} \) shown as black horizontal lines (left-hand panel) for three spectroscopically confirmed LyC leakers and for our two q1 candidates. Since different studies have different strategies of presenting estimated ionizing escape fractions it is important to note that for Q1549–C25 \( f_{\text{esc}} > 51 \) per cent is estimated at 95 per cent confidence, where less than 45 per cent of the \( f_{\text{esc}} \) distribution is \( \leq 100 \) per cent and 95 per cent of the distribution is at \( f_{\text{esc}} > 51 \) per cent (Shapley et al. 2016). In the case of the Ion2 (de Barros et al. 2016; Vanzella et al. 2016) \( f_{\text{esc}} \) values are described in the range from 20 to 100 per cent and for the Ion3 (Vanzella et al. 2018) only \( f_{\text{esc}} \) is reported in range 10 to 100 per cent. Here, we are presenting \( f_{\text{esc}} \) for q1 candidates as a range, where we also including the estimated ranges of the less likely cases of IGM attenuation, \( f_{\text{esc}}^{\text{IGM}} \) and \( (1-D_b) \) and \( \tau_{\text{LyC}}^{\text{IGM}} \), as blue and green lines, respectively. Although we see no clear correlation between EW(Ly\( \alpha \)), \( f_{\text{esc}} \) and \( R_{\text{obs}} \), with the current size of our sample we are not able to rule out its existence. The right-hand panel of the Fig. 10 shows EW(Ly\( \alpha \)) as a function of \( R_{\text{obs}} \). It is interesting to note that the q1 candidates and confirmed detections in the literature in the right-hand panel of Fig. 10 show an absence of strong EW(Ly\( \alpha \)) for higher values of \( R_{\text{obs}} \).

However, with current sample size and the non-uniform selection methods (D10K, VUDS, and 3D-HST) in this work, as well as the fact that EW(Ly\( \alpha \)) line is affected by different processes (morphology, transparency of the IGM, high star formation, etc.), it is not possible to rule out the existence of a correlation or anticorrelation between emitted LyC flux into IGM and properties of the Ly\( \alpha \) line. At this stage, the possible lack of the any clear trend indicates a need for a homogeneous sample selection method and larger sample.

5.3 Position of the LCGs on the colour–colour diagram

The Lyman break technique exploits the expected drop in flux at the Lyman limit (~912 Å) to develop selection criteria on colour–colour diagrams (e.g. Steidel et al. 1996) as an efficient method of selecting 2.5 \( \lesssim z \lesssim 3.5 \) and 3.5 \( \lesssim z \lesssim 4.5 \) SFGs. These LBGs have been used many previous studies searching for LyC emitting galaxies (e.g. Iwata et al. 2009; Vanzella et al. 2010b; Nestor et al. 2011). The limitations of the Lyman break selection in the context of \( z \sim 3-4 \) LCG detection was recognized by Cooke et al. (2014). The authors discuss how the effect of LyC on the \( u \)-band magnitude moves the colours of a galaxy on the colour–colour diagram from their expected LBG location for a given redshift in a predictable manner.

Historically, the templates and composite spectra used to determine the expected positions of LBGs on colour–colour plots and, subsequently, the LBG selection regions, assumed zero LyC flux in order to use the break in flux as a selection discriminant, as a negligible fraction was expected. However, various levels of LyC flux enter the \( u \)-band filter and act to move the colours of the galaxies downward on a \( u-g \) versus \( g-i \) plot, below the locations expected for their redshifts when assuming zero LyC flux, and this effect can be sufficiently strong to place the galaxy colours outside the standard LBG selection region box. For LCGs with redshifts at the lower end of the redshift range being probed, the movement in colour is small (regardless of the level of LyC flux), but the movement increases...
(rather dramatically) toward the higher redshift end, as more LyC flux enters the r band. In fact, for typical $u - g$ versus $g - i$ diagrams, galaxies with roughly $z \gtrsim 3.3$ are positioned well off the diagram (upward in $u - g$ colour, reaching infinity for $z \gtrsim 3.4$). Thus, the very presence of $z \gtrsim 3.3$ galaxies on these plots is indicative of the presence of LyC flux.

As a result, for $z \sim 3$ galaxies, inspecting their locations on conventional $u - g$ versus $g - i$ diagrams is informative regarding the potential for escaping LyC flux. In particular, those galaxies with colours in the region below and outside the $z \sim 3$ LBG selection region (on conventional $u - g$ versus $g - i$ diagrams), as they are expected to have the highest fraction of escaping LyC flux. We note that this region of the diagram (below and outside the conventional selection region) contains the colours of a significant fraction of low-redshift galaxies that makes an efficient selection of $z \sim 3$ LCGs difficult when using broadband optical colours alone.

However, and perhaps more importantly, this same $u - g$ versus $g - i$ diagram provides a good means to select $z \sim 4$ LCGs. As mentioned above, LBGs at $z \gtrsim 3.3$ with zero LyC flux have colours that reside off this diagram, but $z \gtrsim 3.3$ galaxies with even small levels (e.g. $\gtrsim 1$ per cent) of LyC flux are found on this plot and in locations relatively removed from the high density of low-redshift galaxy colours. As a result, $u - g$ versus $g - i$ colour–colour plots, conventionally used for $z \sim 3$ LBGs, are powerful for $z \sim 4$ LCG colour selection, in particular when combined with infrared broadband colours.

The two q1 Lyc candidates from this work are plotted on the $u - g$ versus $g - i$ colour–colour diagram in Fig. 11, where the $z \sim 2.7$–3.4 LBG selection region is shaded in grey. Four dotted teal, green yellow, and red curves are evolutionary tracks of four different LBG composite spectra with different EW(Ly$\alpha$) values adopted from Shapley et al. (2003). Here, all four composite spectra follow the standard LBG colour-selection convention and assume no flux shortward of 912 Å. As we can see from Fig. 11, the dotted line tracks are located inside the LBG selection region and the $u - g$ colour becomes redder as the redshift of the object increases. The locations of the LBGs at $z \gtrsim 3.3$ with zero LyC flux are above the plot and reach infinity by $z \sim 3.4$. Adopting the flat model from Cooke et al. (2014) for all four composite spectra, which provides an average measure of the LyC flux within the $u$ filter, artificial flux is added shortward of 912 Å in different amounts: $R_{\text{obs}} = 1, 2, 5, 10, \text{ and } 20$ per cent. Depending on their redshift and amount of ionizing radiation, we can see how the evolutionary tracks, solid teal, green, yellow, and red lines, indicate that galaxies with the largest fraction of observed LyC flux shift significantly downward in $u - g$ colour, with a fraction of the $z \sim 3$ galaxies residing outside the standard LBG colour-selection region and the new presence of $z \gtrsim 3.3$ galaxies (black triangles and squares indicate $z = 3.5$ and 4.0, respectively). If the assumption from Cooke et al. (2014) is correct then the positions of q1 candidates (blue squares) on the colour–colour diagram are consistent with this prediction and an approximate value of $R_{\text{obs}}$ can be read off the plot.

We also examine positions of the two q1 candidates on the $g - r$ versus $r - i$ colour–colour diagram, right-hand panel of Fig. 11. The $g - r$ versus $r - i$ colour–colour diagram is designed to select LBG galaxies at redshift $3.5 \lesssim z \lesssim 4.5$. Similar to the $u - g$ versus $g - i$ colour–colour plots for selecting $z \sim 3$ LBGs and $z \sim 3$–4 LCGs, the $g - r$ versus $r - i$ colour–colour diagram is useful for selecting $z \sim 4$ LBGs and $z \sim 4$–5 LCGs. As was the case for the $u - g$ versus $g - i$ colour–colour plots, galaxies at the lower redshift end of the redshift range probed (here, $3.5 < z < 4.5$) have little movement in colour, regardless of LyC flux levels, and the plot is more effective in identifying LCGs at the higher end of the redshift range and beyond to $z \sim 5$.

On the right-hand panel in Fig. 11, the dotted evolutionary tracks of the four composite spectra are shown with no LyC flux, while solid lines are evolutionary tracks with added LyC flux. Black triangles and squares denote the expected positions of $z = 4$ and 4.5 galaxy colours, respectively. The evolutionary tracks are colour coded in the same way as in the left-hand panel. Here, we are showing two LBG selection regions. The purple selection region is adopted from Hildebrandt et al. (2009) and designed for CFHT $g, r,$ and $i$ filters that have almost the same bandwidths as Subaru HSC filters. A more inclusive selection region based on the (dotted) evolutionary tracks to $z \sim 4.5$ is shown in grey. Our q1 candidates are marked as blue squares. As a guide, red circles mark the positions of the colours for galaxies with similar redshifts and EW(Ly$\alpha$),...
but without LyC flux. The candidates have bluer colours for their redshifts, as compared to LBGs (i.e. with zero LyC flux), and are consistent with the expectations for galaxies with measurable escaping LyC flux. In a similar manner to the effectiveness of the $z \sim 3$ LBG $u-g$ versus $g-i$ diagram to identify $z \sim 4$ LCGs, this $z \sim 4$ LBG diagram is more effective in identifying LCGs at $z \sim 5$.

### 5.4 Non-detection subsample stacking

For the candidates from the Non-detection subsample, we attempt to estimate the $f_{\text{esc}}$ limit by stacking individual $u$-band images. The data are stacked in two ways, based on their average and median. The stacking procedure is performed in IRAF using the IMCOMBINE task and flux is measured in a 1.2 arcsec circular aperture, the same way as was done with the single galaxy candidates.

Before stacking we inspect the 1D and 2D spectra of all 87 candidates of the Non-detection subsample to ensure $z > 3.42$ redshifts. Two candidates are recognized as AGN and are excluded from further analysis. For the remaining 85 candidates, we confirm redshifts for 39. To create stacks free from flux emitted by low-redshift objects, we only use the 39 Non-detection candidates for which we can spectroscopically confirm redshifts. Along with the stacking of the CLAUDS $u$-band images that probe the LyC UV radiation, we stack $i$, $z$, or $y$ Subaru HSC images for non-ionizing UV flux depending on the redshift of the object. To estimate non-ionizing UV flux for the Non-detection subsample, we used the $i$ band for objects at $3.5 < z < 4.5$, $z$ band for objects at $4.5 < z < 5.5$, and $Y$ band for objects at $5.5 < z < 6.5$. We create two different stacks, average and median, for each redshift bin and the non-ionizing flux is measured the same way as for $u$-band images. This approach results in three average flux values, one for $i$, $z$, and $y$ bands, and three median flux values. Finally, to estimate the average non-ionizing UV flux we averaged the fluxes in $i$, $z$, and $y$-band stacks. The same is done for median values. No LyC detection is seen with S/N $> 3$ in either the average or median stack. The results are presented in Fig. 12.

The estimated magnitude upper limit for the stack of the 39 galaxies is $u_{\text{avg}}^\text{abs} \sim 33.17$ mag and $u_{\text{med}}^\text{abs}$ was negative. In addition, we estimate the upper limits for the averaged flux density ratio $R_{\text{obs}}$ and the averaged $f_{\text{esc}}^\text{abs}$ to be 0.001 and 0.006, respectively. For the intrinsic luminosity ratio we adopt $\langle L_{\text{LyC}}/L_{\text{1500}} \rangle_{\text{obs}} = 0.3$. The $E(B-V)$ values for the 39 Non-detection galaxies (with confirmed redshifts) are extracted from the catalogue published by Laigle et al. (2016). Based on the 30 objects cross-matched by their coordinates, we adopt the average $E(B-V) = 0.084$ for the 39 spectroscopically confirmed Non-detection candidates. It has been estimated that, if galaxies alone are sufficient to reionize the Universe, we would expect them to emit $\sim 10 - 20$ per cent of their LyC radiation on average into the IGM (Robertson et al. 2015; Khaire et al. 2016). The estimated upper limits on $f_{\text{esc}}^\text{abs}$ from the average and median stacks indicate that the emission of LyC radiation into IGM from galaxies is not greater than 1 per cent. From the results estimated...
by stacking, and their associated assumptions, we can conclude that
the average amount of LyC radiation that escapes into IGM from
the galaxies at 3.5 < z < 5.1, if representative of galaxies at z > 6,
is not enough to sufficiently support reionization.

6 SUMMARY AND CONCLUSIONS

In this work, we present the results of a search for a population
of LCGs whose progenitors at z > 6 may be responsible for
reionizing the Universe. We used deep CLAUDS u-band photometry
in the COSMOS field to search for LyC flux emitted by LCGs. To
ensure the cleanest sample, and that our measured flux was not
contaminated by low-redshift interlopers, we only include objects
with high-quality spectroscopic redshifts from the D10K, VUDS,
and 3D-HST surveys. HST and ground-based imaging in multiple
filters were examined individually for every object that potentially
showed LyC flux in the u band and the 1D and 2D spectra were
analysed to help eliminate low-redshift contaminants next to the
galaxies in the line of sight. After the selection process, we identify
five candidates within the redshift interval 3.5 < z < 5.1, which we
divide into two quality groups q1 (2 objects) and q2 (3 objects). Our
main conclusions can be summarized as follows:

(i) Following the approach of Cooke et al. (2014), we use the z ∼
4 LCG u − g versus g − i colour–colour diagram (conventionally used
for z ∼ 3 LBG selection) to investigate our candidates. The positions
of the q1 LCG candidates on the u − g versus g − i colour–colour
diagram are consistent with galaxies having measurable escaping
LyC flux (here, R_{obs} ∼ 15 per cent) for galaxies with their redshifts
and EW(Lyα). We note that R_{obs} is not f_{esc}.

(ii) Adopting the mean IGM transmission values as representa-
tive for single LCG candidates leads to overestimated f_{esc} values
and indicates that any detection of LyC at z > 3.5 is likely to
arise from a more transparent than average sightlines. On the other
hand, IGM corrections based on maximum transparency to LyC
photons give underestimated f_{esc} values. These results imply
that f_{esc} measurements for single objects at high redshifts can only
be determined with broad range of a values and the only way for
improvement is to better understand impact of the IGM on a single
line of sight.

(iii) The estimated f_{esc} are in the range ∼ 5 − 73 per cent for q1
candidate ID 1 and ∼ 30 − 93 per cent for q1 candidate ID 368,
where different metallicities can produce a variation in f_{esc} up to
∼ 10 per cent.

(iv) Both q1 LCG candidates have EW(Lyα) < 50 Å. No clear
correlation is seen among EW(Lyα), f_{esc} and R_{obs} in our relatively
small sample. We note that we do not see cases where an LCG
candidate has strong EW(Lyα) and high R_{obs} ratio. With current
size of the sample we are unable to exclude the existence of a
correlation or anticorrelation between EW(Lyα) and LyC leakage
into IGM.

(v) The stacking procedure of the Non-detection candidates did
not reveal any significant LyC flux above S/N > 3 in the u band.
Based on the results from stacking 39 Non-detection candidates
with confirmed redshifts, the LyC radiation emitted by galaxies into
IGM does not exceed 1 per cent for the average stack. If this is the
case for galaxies at z > 6 then galaxies alone are not able to emit
enough LyC radiation to reionize the Universe.

Our analysis demonstrates that the u − g versus g − i colour–
colour diagram is useful for identifying z ∼ 4 LCGs. However, it
emphasizes that detecting clean, reliable sources of LyC radiation
is difficult without using high-resolution HST imaging in at least
one filter, and preferably two or more, as well as spectroscopic
redshift confirmation in combination with deep ground- or space-
based LyC photometry. Creating larger samples of LCGs at z > 3
will be possible with large spectroscopic surveys that are
followed by high spatial resolution imaging with new ground-
based 30-m telescopes and space-based telescopes such as the Large
UV/Optical/IR Surveyor (Bolcar et al. 2017) or the Cosmological
Advanced Survey Telescope for Optical and UV Research (Côte
et al. 2012).

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APPENDIX A: IMAGES OF THE Q1 AND Q2 CANDIDATES IN OTHER AVAILABLE FILTERS

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Figure A1. Cutouts of the q1 (ID 368 and 1) and q2 (ID 326, 330, and 421) candidates in g-HSC, r-HSC, i-HSC, and z-HSC as well as HST F125W, F140W, 160W, and F606W where available. The green circles are 2 arcsec in diameter and thumbnails are 15 arcsec × 15 arcsec in size.
APPENDIX B: SPECTRA OF THE Q1 AND Q2 CANDIDATES
Figure B1. 1D spectrum of the selected five LCG candidates (q1 and q2). Composite spectrum (blue) is fitted to the candidate spectrum (grey) to confirm reported redshift, detect Lyα forest or other spectroscopic features (blue dashed lines). With yellow vertical stripes positions of the sky lines are marked.
Figure B2. Same as Fig. B1.
Figure B3. Same as Fig. B1.

Figure B4. Keck MOSFIRE 1D and 2D spectrum of the candidate ID 368 used to additionally confirm redshift of the candidate. Positions of the $\text{H}$\textbeta\ and [O\textsc{iii}] lines are marked.

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