The 2175 Å bump features in FeLoBAL quasars: One indicator of MW-like dust in the nuclear region of quasars

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ABSTRACT

To investigate the properties of dust in the nuclear region of quasars, we explored the extinction curves of the iron low-ionisation broad absorption line (FeLoBAL) quasar SDSS J163004.29+311957.6 and its two analogues. The parametrised extinction curves indicated the Milky Way-like 2175 Å bump features in underlying extinction, which are similar to those seen in the Local Group and in a subset of high-redshift star-forming galaxies. Compared to the bump features in the Large Magellanic Cloud (LMC), the detections in this work are much closer to those in the Milky Way (MW). These bump features, as well as those in the high- and low-ionisation broad absorption line (BAL) quasars with 2175 Å bumps, are probably the counterpart of the 2175 Å bump features in the quasar environment. This type of dust grain is generally small, easily disrupted by high-energy photons, and has difficulty surviving in the radiation field of the active galactic nucleus. However, due to the presence of absorption-line outflows, the 2175 Å bump feature in quasars, which should be rare, is seen many times in BAL quasars. The shielding effect of outflow clouds allows the MW-like dust grains to be assembled or extends the survival period in the quasar nuclear region. The process and physical and chemical conditions deserve further observational study and investigation.

Key words. dust, extinction – galaxies: ISM – quasars: absorption lines – quasars: individual: SDSS J163004.29+311957.6

1. Introduction

Dust is everywhere in the universe, and although this dust is tiny, it plays an important role in many astrophysics processes. Absorption, scattering, and re-radiation of the dust reshape the spectral energy distribution (SED) of galaxies and active galactic nuclei (AGNs). At the same time, as an important component of the interstellar medium (ISM), the dust grain surface offers a place to efficiently form molecules as a critical step for star formation, which creates galaxies (Draine 2003). Understanding dust properties and their evolution is crucial to providing complete knowledge of ISM processes, star formation, and galaxy evolution.

In the galaxy, dust fills interstellar space. Its distribution is heterogeneous, and it is known to be clumpy with cold, dense regions of star formation. In the AGN, the dust generally accumulates at the periphery of the accretion disk and forms a circumnuclear dusty torus as an important structural element in the AGN unified schemes (e.g. Antonucci 1993; Urry & Padovani 1995). The torus obscures and absorbs high-energy photons emitted from the central engine and re-radiates at thermal infrared (IR) wavelengths (e.g. Rees et al. 1969; Barvainis 1987). It is therefore also considered to be the reason for the appearance of type I/II AGNs. Beyond this, a large amount of dust is still distributed in various parts of the nuclear region, such as AGN outflows. Studies have found that high-speed outflows have an unexpectedly significant correlation with hot dust emission, which suggests a dust outflow scenario (e.g. Wang et al. 2013; Zhang et al. 2014).

Dust attenuation, polarimetry, and thermal radiation are all efficient observational means for dust. The blocking effect on centre photons implies the existence of a dusty torus, and polarimetric and IR spectroscopic or photometric observations have presented abundant information about the physical and structural characteristics (e.g. Krolik & Begelman 1988; Pier & Krolik 1992; Dullemmond & van Bemmel 2005; Tristram et al. 2007, 2014; Nenkova et al. 2008; Alonso-Herrero et al. 2011; García-González et al. 2017; Zhuang et al. 2018; Lopez-Rodriguez et al. 2020; Lyu & Rieke 2021). Relatively speaking, dust attenuation can delve into the details of interstellar dust grains, such as the size distribution of dust grains (e.g. Cardelli et al. 1989; Whittet 2003) and even the chemical compositions (e.g. Li & Draine 2001; Draine 2003). Furthermore, measuring the extinction curve and making a corresponding reddening correction allows accurately deriving the intrinsic SED, luminosity functions, and physical properties of AGNs.

Regarding the AGN extinction law, Czerny et al. (2004) compared the composites of red and blue quasars, and Gaskell et al. (2004) also studied the composites of quasars with...
different radio properties. These composite comparative studies found that quasar extinction curves are significantly flatter in ultraviolet (UV) than those of local galaxies on average, and they do not show any traces of the 2175 Å bump feature, which is ubiquitous in extinction curves of the Milky Way (MW) and Large Magellanic Cloud (LMC) (see Draine 1989). Similar results have been obtained from the analyses of the quasar colour distribution (Richards et al. 2003; Hopkins et al. 2004), and the red tail of the colour distribution of quasars is well described by Small Magmatic Cloud (SMC) – like reddening. However, Jiang et al. (2013) reported a quasar intrinsic anomalously steep reddening law via the exceptional source IRAS 14026+4341, and the extinction curves detected in all the individuals of the high Av quasar sample are also steeper than the SMC law (Zafar et al. 2015). In the study of polar dust obscuration in nearby broad-line active galaxies (BLAGNs) from the XMM-XXL field, Buat et al. (2021) reported that even 60% of the sources containing polar dust exhibit a steep extinction curve. The reality is probably more complex, with a mix of different shapes (e.g. Buat et al. 2021), and we are still far from understanding quasar extinction phenomena. Moreover, Zhang et al. (2017) derived the SMC-type reddening law of a dusty torus in three narrow-line Seyfert 1 galaxies (NLS1s). However, the extinction-to-gas ratios \((E(B-V)/N_{H})\) of the two cases are comparable to the Galactic standard value (Bohlin et al. 1978) and approximately 10 to 20 times higher than that of the SMC (Weingartner & Draine 2001). These findings indicate the dominance of small dust grains in the quasar environment, which offers the possibility of detecting the 2175 Å bump feature in quasars.

To date, dozens of 2175 Å bump features have been reported at high redshifts. These detections are mostly found in individual intervening systems towards background quasars, including metal (e.g. Mg II, Zn II, and C I) absorbers (e.g. Wang et al. 2004; Srianand et al. 2008; Noterdaeme et al. 2009; Zhou et al. 2010; Jiang et al. 2010, 2011; Ledoux et al. 2015; Ma et al. 2017) and DLAs (e.g. Wucknitz et al. 2003; Junkkari et al. 2004; Prochaska et al. 2009; Wang et al. 2012; Ma et al. 2015), or towards gamma ray burst (GRB) afterglows (e.g. Ellison et al. 2006; Vreeswijk et al. 2006; Prochaska et al. 2009; Liang & Li 2010; Schady et al. 2012; Zafar et al. 2011, 2012, 2018). These systems have a high metallicity and high dust depletion level (e.g. Ma et al. 2017). These detections reveal information about the ISM in the foreground galaxies of quasars or GRBs or GRB hosts rather than the quasar itself. The quasar-associated 2175 Å bump features are still rarely detected, and they have been found in only a limited number of quasar intrinsic extinction curves (Zhang et al. 2015a; Pan et al. 2017; Shi et al. 2020). Theoretically, the carriers of 2175 Å bump features are small carbonaceous and silicate grains (e.g. Joblin et al. 1992; Weingartner & Draine 2001; Li & Draine 2001; Steglich et al. 2010; Blasberger et al. 2017); however, they are easily disrupted by high-energy photons, and they can hardly exist in the radiation field of the centre engine (e.g. Voit 1992 and reference therein). The question is why this type of dust grain can be produced and survive in this harsh environment. This is an interesting topic worth of further investigation. Exploring individuals, sample expansion, and further detailed studies of important targets may provide a possibility for researching the formation and destruction of 2175 Å dust grains in the nuclear region of quasars.

Broad absorption lines (BALs) are the most famous quasar intrinsic absorption systems, and they are one of the direct manifestations of AGN outflows. They appear in the spectra of ~10–26% of optically selected quasars and are often observed as absorption by the ions of C IV, Si IV, Al III, and Mg II (e.g. Weymann et al. 1991; Reichard et al. 2003; Trump 2006; Gibson et al. 2009; Scaringi et al. 2009; Zhang et al. 2010; Allen et al. 2011). Moreover, iron low-ionisation broad absorption lines (FeLoBALs) are a rare and special subclass of BALs, which exhibit absorption from excited fine-structure levels or excited atomic terms of Fe II or Fe III (e.g. Hazard et al. 1987; Becker et al. 1997, 2000; Menou et al. 2001; Hall et al. 2002; Zhang et al. 2015b). The mid-IR spectroscopic observation of six FeLoBAL quasars detected unidentified emission signatures of carbonaceous molecules (Farrah et al. 2010), which implies that the carriers of 2175 Å bump features can probably survive in the environments of BAL quasars. In Zhang et al. (2015a), the initial exploration of 2175 Å bump features among Mg II absorption line systems on quasar spectra in the SDSS DR10 led to the identification of excess broadband extinction near 2250 Å in seven BAL quasars. These cases are in the redshift range of 1.78 to 2.29, and five of them show low-ionisation BALs (LoBALs) of Mg II and Al III. This provides direct evidence of the coexistence of BALs and 2175 Å bump features in quasars.

In this work, we report the detection of the 2175 Å bump feature in three FeLoBAL quasars. The remainder of this paper is organised as follows: in Sect. 2 we describe the observational data available and data reduction. In Sect. 3 we present the derivation of the extinction curve, Sect. 4 provides a discussion, and Sect. 5 summarises the main conclusions.

## 2. Observations and data reduction

After Zhang et al. (2015a), 2175 Å bump features in the spectra of the Sloan Digital Sky Survey (SDSS; York & Adelman 2000) DR7 FeLoBAL quasars (Shen & Richards 2011) were systematically explored following the search procedure of Zhang et al. (2015a). The quasar composite reddened by a parametrised extinction curve (introduced in the first two paragraphs of Sect. 3) was used to fit the observed spectra in the quasar rest frame and obtain potential bump candidates, and then a simulation technique of the control sample (Jiang et al. 2010; also see the fifth paragraph of Sect. 3) was performed to gauge the significance of the bump features. Ten bump candidates were screened out with 2175 Å bump features at a >3σ level of statistical significance (see Table 1). After performing a visual examination, seven false signals were rejected because their spectral wavelength ranges only cover part of the bump feature, and no effective photometric data limit the far-UV models. We finally identified three credible candidates, that is, J102036.10+602339.0 (referred to as SDSS J1020+6023 hereafter), SDSS J134951.93+382334.1 (referred to as SDSS J1349+3823 hereafter), and SDSS J163004.29+311957.6 (referred to as SDSS J1630+3119 hereafter). Their emission-line redshifts given by the SDSS DR7 quasar catalogue (Schneider et al. 2010) are 0.9940 ± 0.0006, 1.0943 ± 0.0009, and 1.9960 ± 0.0007, respectively. Since SDSS J1630+3119 has the most observational data and its spectrum shows the entire bump profile, it was analysed and is reported in the text in detail. For SDSS J1020+6023 and SDSS J1349+3823, their optical spectra cannot completely cover the bump profiles because of their lower redshifts; fortunately, the reliability of spectral fitting is verified by the GALEX data; thus, they are listed in the appendices.
Table 1. Best-fit parameters of ten bump candidates in FeLoBALQs.

| Name (SDSS J)   | z    | c1     | c2     | x0    | γ    | Significance |
|-----------------|------|--------|--------|-------|------|-------------|
| 101927.37+022521.4 | 1.3643 | −1.95 ± 0.02 | 0.65 ± 0.01 | 0.76 ± 0.06 | 4.46 ± 0.01 | 0.99 ± 0.04 | 5.6σ |
| 102036.10+062339.0 | 0.9940 | −1.82 ± 0.02 | 0.54 ± 0.01 | 0.95 ± 0.08 | 4.49 ± 0.01 | 1.20 ± 0.03 | 6.9σ |
| 102249.27+130125.0 | 1.2168 | −0.87 ± 0.03 | 0.55 ± 0.01 | 1.41 ± 0.07 | 4.78 ± 0.02 | 1.68 ± 0.04 | 3.7σ |
| 110711.40+082331.2 | 1.3875 | −1.73 ± 0.04 | 0.81 ± 0.01 | 0.57 ± 0.04 | 4.69 ± 0.01 | 0.95 ± 0.05 | 6.2σ |
| 112220.76+153927.8 | 1.0996 | −1.38 ± 0.01 | 0.43 ± 0.01 | 1.11 ± 0.06 | 4.56 ± 0.01 | 1.10 ± 0.02 | 7.0σ |
| 134951.93+382334.1 | 1.0943 | −0.73 ± 0.02 | 0.37 ± 0.01 | 2.02 ± 0.13 | 4.67 ± 0.01 | 1.18 ± 0.08 | 17.2σ |
| 142010.28+604722.3 | 1.3450 | −0.58 ± 0.03 | 0.51 ± 0.01 | 0.39 ± 0.06 | 4.58 ± 0.01 | 0.87 ± 0.04 | 8.7σ |
| 152438.79+415543.0 | 1.2299 | −1.34 ± 0.01 | 0.41 ± 0.01 | 0.70 ± 0.06 | 4.47 ± 0.01 | 0.96 ± 0.03 | 10.9σ |
| 155633.78+351757.3 | 1.4950 | −3.46 ± 0.02 | 1.00 ± 0.01 | 1.35 ± 0.10 | 4.56 ± 0.01 | 1.45 ± 0.02 | 7.6σ |
| 163004.29+311937.6 | 1.9960 | −3.37 ± 0.05 | 1.11 ± 0.01 | 0.90 ± 0.07 | 4.51 ± 0.01 | 0.92 ± 0.07 | 31.2σ |

Table 2. Broad-band photometric data of SDSS J1630+3119.

| Band | Magnitude | Date-observation | Survey |
|------|-----------|------------------|--------|
| u'   | 23.15 ± 0.46 | 04/29/2003       | SDSS   |
| g'   | 21.41 ± 0.04 | 04/29/2003       | SDSS   |
| r'   | 20.12 ± 0.02 | 04/29/2003       | SDSS   |
| i'   | 19.16 ± 0.02 | 04/29/2003       | SDSS   |
| z'   | 18.31 ± 0.03 | 04/29/2003       | SDSS   |
| J    | 17.420 ± 0.221 | 04/05/1998      | 2MASS  |
| H    | 16.178 ± 0.170 | 04/05/1998      | 2MASS  |
| Ks   | 15.871 ± 0.205 | 04/05/1998      | 2MASS  |
| W1   | 15.217 ± 0.205 | 02/17/2010      | WISE   |
| W2   | 14.080 ± 0.044 | 02/17/2010      | WISE   |
| W3   | 10.720 ± 0.076 | 02/17/2010      | WISE   |
| W4   | 8.191 ± 0.206  | 02/17/2010      | WISE   |
| V    | –           | 07/06/2005–10/22/2013 | CSS   |
| g    | –           | 06/17/2009–06/27/2013 | Pan-STARRS |
| r    | –           | 06/07/2010–06/20/2013 | Pan-STARRS |
| i    | –           | 06/09/2009–05/11/2014 | Pan-STARRS |
| z    | –           | 03/01/2010–08/16/2013 | Pan-STARRS |
| y    | –           | 02/20/2010–08/22/2013 | Pan-STARRS |
| zr   | –           | 06/11/2018–09/02/2019 | ZTF   |

SDSS J1630+3119 is an IR-luminous quasar with an extremely faint UV emission. The broad-band SED for SDSS J1630+3119 from UV to near-IR is presented by the photometric images taken with the SDSS at the u', g', r', i', and z' bands, the Two Micron All Sky Survey (2MASS; Skrutskie & Cutri 2006) at the J, H, and Ks bands, and the Wide-field Infrared Survey Explorer (WISE; Wright & Eisenhardt 2010) at the W1, W2, W3, and W4 bands. The optical and near-IR photometric data of SDSS J1630+3119 join smoothly, suggesting that the quasar did not vary significantly between the SDSS and 2MASS observations. The multi-band magnitudes are summarised in Table 2.

SDSS J1630+3119 was targeted for spectroscopic observations by the SDSS as a high-redshift quasar candidate (LEGACY_TARGET1 = QSO_HIZ). Two archival observations for SDSS J1630+3119 from the SDSS and SDSS-III/BOSS were performed on 4 April 2005 (DR7; Abazajian et al. 2009) and 7 July 2011 (DR16; Dawson & Kneib 2016), and the 1D spectra were accessed from the SDSS Science Archive Server. The BOSS spectrum multiplied by approximately 1.2 times is in excellent agreement with the SDSS spectrum. The Catalina Sky Surveys (CSS; Drake et al. 2009), the Panoramic Survey Telescope & Rapid Response System (Pan-STARRS; Chambers et al. 2016), and the Zwicky Transient Facility (ZTF; Masci et al. 2018) show that SDSS J1630+3119 does not have significant long-term variability in either optical or near-IR bands within the current measurement errors for 14 years beginning on 6 July 2005 (Fig. 1). Thus, the spectral flux density differences between the two spectroscopic observations probably come from the spectrophotometric calibration of the different SDSS optical spectroscopy pipelines (Version ‘26’ and ‘v5_13_0’). Moreover, the scaled BOSS spectrum agrees very well with the SDSS photometric data, which indicates that the recalibration is reliable and that SDSS J1630+3119 did not vary significantly between the SDSS photometric and spectroscopic observations. In the following analysis, the scaled BOSS spectrum is used to derive the extinction curve because its wavelength coverage is wider. All of the data were corrected for the Galactic reddening of $E(B-V) = 0.032$ (Schlegel et al. 1998) and the Fitzpatrick (1999) reddening curve before any further analysis.

The spectrum of SDSS J1630+3119 shows unusual absorption features and has been classified as a BAL quasar in the optical by Masci et al. (2018). The broad-band photo-
by 0.75 to match that of SDSS J1630+3119 and BAL troughs. The spectrum was slightly smoothed with a Gaussian kernel with a sigma of 3 spectral pixels to help show the absorption features. The spectrum of SDSS J1630+3119 is dominated by strong broad and blueshifted absorption troughs from Mg II, Al III, C IV, and Fe II UV 1, UV 2, UV 3, UV 62.63, Opt. 6,7 and Opt. 8 multiplet transitions, the vertical dash-dotted grey marks indicate the expected rest-frame positions of the resonant lines in these features. The comparison spectrum (shown in green) with similar absorption troughs to SDSS J1630+3119 from the FeLoBAL quasar SDSS J1440+3710 was smoothed and multiplied by 0.75 to match that of SDSS J1630+3119 in the wavelength ranges of 1600–1900 Å and 2800–3300 Å.

Gibson et al. (2009) and Shen & Richards (2011). In Fig. 2 we present the observed spectrum of SDSS J1630+3119 and the absorption troughs of C IV, Al III, Mg II, and iron multiplets (marked by dot-dashed grey lines and characters). A typical FeLoBAL quasar with absorption troughs similar to those of SDSS J1630+3119 and SDSS J144002.24+371058.5 (referred to as SDSS J1440+3710 hereafter) is over-plotted for comparison, in which the spectrum is multiplied by a factor to match that of SDSS J1630+3119 in the wavelength ranges of 1600–1900 Å and 2800–3300 Å. The high-ionisation absorption lines of C IV λ1548, 1551 in the spectrum of SDSS J1630+3119 are revealed by troughs near the centre of the corresponding lines, which have an extremely broad velocity width from approximately 4000 to ~13 000 km s⁻¹ with respect to the systemic velocity. The numerous absorption lines from low-ionisation species, such as Mg II, Al III, and Fe II, that characterise the spectrum show some different appearances that are narrower and heavily blended together for most of them in the lower-resolution spectrum. In the figure, we identify some strong absorption troughs of Fe II multiplets UV 1, UV 2, UV 3, and UV 62.63 near 2600 Å, 2382 Å, 2344 Å, and 2750 Å (in the quasar rest frame), and other weak multiplets of Fe II UV 8 and UV 44.45 longward of the C IV emission line, and Fe II UV 8, UV 145, and UV 263 on either side of the UV 1 multiplets are not marked.

3. Extinction curve and bump

In Fig. 3a we present the observed data of SDSS J1630+3119, including the original spectrum (black curve) and multi-band magnitudes (orange squares). The dramatic flux drop shortward of the Mg II broad emission line (BEL) appears to slow down somewhere between the SDSS r' and i' bands, suggesting that a broad absorption bump occurs there. In the MW and other galaxies in the Local Group, the interstellar extinction is most commonly obtained through the pair method by comparing the spectra of two stars of the same spectral type, one of which is reddened (f_obs), while the other is unreddened (f_model). A similar technique, called the quasar spectrum pair method, was used to derive the quasar extinction curve and identify 2175 Å bump features in the quasar spectrum (e.g. Wang et al. 2004; Zhou et al. 2010). The reddened spectrum is an observed quasar spectrum, and the unreddened spectrum is replaced by the quasar composite or an individual blue quasar spectrum. In this work, the unreddened spectrum was set to a quasar composite, which was obtained by combining the SDSS DR7 quasar composite (λ < 3000 Å; Jiang et al. 2011) and the near-IR quasar template (λ ≥ 3000 Å; Glikman et al. 2006).

We used the Fitzpatrick & Massa (1990) parametrisation of the optical/UV extinction curve; however, the extinction curve we measured in this work is A(λ) rather than k(λ − V) (≡ E(λ − V)/E(B − V), used in the original Fitzpatrick & Massa formula). The reason for this is explained in the next paragraph. The above two parameters are equivalent for a given extinction curve. A(λ) is described as a combination of a linear background
and a Drude component, representing the underlying extinction and the possible 2175 Å bump feature, respectively,

\[
A(x) = -2.5 \log \frac{f_{\text{obs}}}{f_{\text{model}}} = c_1 + c_2 x + c_3 \frac{x^2}{(x^2 - x_0^2)^2 + x^2 y^2},
\]

(1)

where \(x\) (\(\equiv \lambda^{-1}\)) is the inverse wavelength and \(x_0\) and \(\gamma\) are the peak position and full width at half maximum (FWHM) of the Drude profile, respectively. The UV linear component is set by the slope \(c_2\) and intercept \(c_1\), which depend on the relative intensity of both model and observed spectra. In this work, the model and observed spectra are not normalized. If they are normalized at the IR wavebands where the extinction is negligible, the parameter \(c_1\) will be zero. The far-UV curvature term (\(x \geq 5.9 \text{ mm}^{-1}\)) in the Fitzpatrick & Massa formula was excluded. The limited wavelength coverages and the low signal-to-noise ratio (S/N) blueward spectra of SDSS J1630+3119 are insufficient for properly identifying the extinction in the far-UV band. Here, the 2175 Å bump feature is thought to have the same redshift as the quasar. If the devised peak position is much larger than the mean value of \(x_0 = 4.59 \pm 0.027 \text{ mm}^{-1}\) in galactic interstellar curves, then the bump in the quasar spectrum would be the intervening system, which is not associated with the quasar. After masking the wavelength range with Mg II, Al III, C IV, and strong iron absorption troughs, the spectrum was fitted with models as the quasar composite that is reddened by dust with the parametrised extinction curve. We performed the least squares minimisation using the Interactive Data Language (IDL) procedure MPFIT developed by Markwardt (2009). The best-fitting model yields \(c_1 = -3.37 \pm 0.05\), \(c_2 = 1.11 \pm 0.01\), \(c_3 = 0.90 \pm 0.07\), \(x_0 = 4.51 \pm 0.01 \text{ mm}^{-1}\), and \(\gamma = 0.92 \pm 0.07 \text{ mm}^{-1}\). The uncertainties are the formal MPFIT errors.

In general, the extinction curves are exhibited in their native normalisation \((E(\lambda - V)/E(B - V))\) and \((E(\lambda)/E(B - V))\) (Cardelli et al. 1989). The parameter \(R_V\) can also serve as an indicator of galactic environmental characteristics, and its value in the dense region is higher than that on the diffuse ISM (Whittet 2003). In this work, the derived curve was not normalised because the optical and UV SEDs of quasars are directly related to the fundamental plane of black hole activity, and the continuum slope is very different in different quasars. We do not know the unreddened intrinsic continuum for a given quasar, and it was set to a quasar composite in the fitting. Obviously, there must be some deviation between them, and it will be reflected in the parameter \(c_2\). Thus, in the extinction measure process of this work, we cannot actually precisely extract the conventionally defined extinction parameters, such as \(A_V\), \(E(B - V)\) and \(R_V\), from the derived extinction curve. Fortunately, the peak position and width of the bump feature are not correlated with one another (Fitzpatrick & Massa 1986), but both are correlated with the slope of the extinction curve (Cardelli et al. 1989). Thus, the detection of the 2175 Å bump feature from the derived extinction curve is viable, and we can investigate the extinction properties of these special dust grains. The bump strength can be defined using these parameters: \(A_{\text{bump}} = \pi c_3 / 2y = 1.53 \pm 0.16 \text{ mm}^{-1}\), which can be interpreted as rescaling the integrated apparent optical depth of the bump feature \((A_{\lambda} \equiv 2 \pi m_\text{g} \tau_{\lambda})\). Likewise, \(A_{\text{bump}}\) is not normalised either. In Gordon et al. (2003) and Fitzpatrick & Massa (2007), the bump strength was defined as the normalised strength by the color excess \((A_{\text{bump}} / E(B - V))\).

In Fig. 3a, the best-fitting model for SDSS J1630+3119 is overplotted with the observed spectra using a solid green curve. To emphasise the requirement of the 2175 Å bump feature in the extinction curve, we also overplot the reddened quasar composite with only the linear component of the best-fitting model (the dashed green curve). The difference between them corresponds to the bump extinction feature (sky blue shadow). In Fig. 3b, we present the observed (orange squares and black curve) and modelled (green curve) extinction of SDSS J1630+3119. The dashed grey and solid blue lines present the underlying and bump components of the best-fitting extinction curve. The extinction curve rapidly decreases with an increasing wavelength into the infrared spectrum, and \(E(x)\) is an unlimited approach to zero with the a wavelength of \(\lambda \rightarrow \infty\).

To confirm the existence of the 2175 Å bump feature in SDSS J1630+3119, we employed several independent validation tests. The first was the simulation technique of the control sample developed by Jiang et al. (2010). This technique begins with the construction of a control sample, including 200 quasars with a minimum redshift differential relative to SDSS J1630+3119 and an i'-band spectral \(S/N > 6\) from the SDSS DR7 quasar catalogue. Then, we fit each of them by reddening a quasar composite with a parametrised extinction curve. The parameters \((x_0, \gamma)\) were fixed to the best-fitting values of the 2175 Å bump feature in SDSS J1630+3119. Because the pseudo-continuum (the complex of nuclear continuum and Fe II emission multiplets) diversity between the quasar spectra and the composite can mimic a false bump feature (e.g. Pitman et al. 2000), the bump strength distribution of the control sample can be considered as the measurement fluctuation of the 2175 Å bump feature. The distribution is expected to be a Gaussian profile, as shown by the pink curve in Fig. 4. In principle, a bump candidate with a significance level of >3\(\sigma\) is considered a real detection, where \(\sigma\) is the standard deviation of the Gaussian profile. The bump strength in SDSS J1630+3119 is obviously far from the distribution \((A_{\text{bump}} = -0.01 \text{ mm}^{-1}\) and \(\sigma = 0.05 \text{ mm}^{-1}\)), and the bump detection in SDSS J1630+3119 has an extremely high statistical significance.

The second test relies more on intuition. It focuses only on the bump features we are interested in. For the purpose of comparison, we again used the typical FeLoBALquasar, SDSS J1440+3710. In Fig. 5 we present the corrected spectrum of
4. Discussion

SDSS J1630+3119 is not the only special case; in the appendices, we also present the analysis of 2175 Å bump features observed towards two FeLoBAL quasars, SDSS J1020+6023 and SDSS J1349+3823. The parametrised extinction curves yield prominent bump features at 2175 Å in the underlying extinction. The bump parameters are \( x_0 = 4.49 \pm 0.01 \, \mu \text{m}^{-1} \), \( \gamma = 1.20 \pm 0.03 \, \mu \text{m}^{-1} \), and \( A_{\text{bump}} = 1.24 \pm 0.11 \, \mu \text{m}^{-1} \) for SDSS J1020+6023 and \( x_0 = 4.67 \pm 0.01 \, \mu \text{m}^{-1} \), \( \gamma = 1.18 \pm 0.08 \, \mu \text{m}^{-1} \), and \( A_{\text{bump}} = 2.68 \pm 0.25 \, \mu \text{m}^{-1} \) for SDSS J1349+3823. Independent tests imply that both detections have extremely high statistical significance levels. Compared with SDSS J1630+3119, the two cases appear to be the same, showing similar absorption troughs of the Fe II multiplets, extreme reddening, and prominent bump features; thus, they are listed as analogues of SDSS J1630+3119 in the appendices. As Figs. A.1 and B.1 show, unlike SDSS J1630+3119, that the spectra of SDSS J1020+6023 and SDSS J1349+3823 cover most of the bump profiles, except for the blueward profile. The GALEX NUV magnitudes are high levels consistent with the reddened quasar composite, which suggests that the extinction shape and bump are well defined.

4.1. Possible associated or intrinsic quasar systems

Usually, absorbers present in quasar spectra are classified as intervening or associated systems based on the blueshift velocity relative to the quasar rest frame. The former is produced by intergalactic clouds or galaxies located far from the background quasars, and the latter originates from materials associated with the host or intrinsically in quasar. The 2175 Å bump features in SDSS J1630+3119 and its analogues have the same redshifts as quasars, which indicates that these features are bound to be quasar-associated.

However, the carriers of these 2175 Å bump features are still unable to be distinguished, whether the star formation region of the host or the nuclear region of the quasar exists. Although some works suggest that the profile and shift of the concomitant metal absorption lines might be able to be a discriminating pointer (e.g. Shi et al. 2020), this undoubtedly gives a higher requirement for spectroscopic observations. In this work, the spectral resolution is far from achieving the degree of resolving line profiles in absorption troughs; meanwhile, there is no narrow emission line in the wavelength coverage to measure the precise system redshift. We might be able to try to explore the origin of bump features by means of the simultaneous occurrence of BALs. On the one hand, the host of quasars may experience strong high-energy radiation from the nucleus, while the quasar feedback on hosts may also produce a disturbance in the ISM. Hence, the probability of 2175 Å bump features existing in hosts is relatively small compared to quiescent galaxies, which is why quasar-associated 2175 Å bump features are rarely found. On
the other hand, previous statistical works indicate that approximately 15% of optically selected quasars show BAL features, and the detection probabilities of LoBALs and FeLoBALs are two and three powers of the above number (e.g. Weymann et al. 1991; Reichard et al. 2003; Trump 2006; Gibson et al. 2009; Zhang et al. 2010). However, the cases found in the literature (Zhang et al. 2015a; Pan et al. 2017; Shi et al. 2020) and this work give a different conclusion. From these works, we collected 28 sources whose 2175 Å bump features have the same redshifts as quasars. Thus, the BAL fractions are nearly 36% (BALs), 29% (LoBALs), and 11% (FeLoBALs) in this small sample. This situation suggests that there is a strong correlation between occurrences of BALs and 2175 Å bump features, and that the 2175 Å bump features we observed are intrinsic to quasars rather than the hosts.

Admittedly, this small sample has a strong selection effect. The exploration in Zhang et al. (2015a) was based on the SDSS DR10 Mg II absorption-line sample, in which the quasars with unusual absorption lines (i.e. FeLoBALs) were removed through visual inspection. However, this work begins from the known SDSS DR7 FeLoBAL quasars; therefore, the fraction of the FeLoBAL-2175 Å bump features is relatively grossly underestimated. Moreover, statistical studies of optically selected BAL quasars indicate that LoBAL quasars show a moderate level of reddening and FeLoBAL quasars are heavily reddened (e.g. Hall et al. 2002; Reichard et al. 2003; Zhang et al. 2010; Dunn et al. 2015), and the additional 2175 Å bump features make these objects fainter and easier to miss in the magnitude-limited sample. If an unbiased systematic exploration for 2175 Å bump features is performed in homogeneous quasar samples, the foregoing conclusion will probably be greatly enhanced.

### 4.2. Comparison of the bump parameters

Figure 6 illustrates the bump parameters of the 2175 Å bump features in SDSS J1630+3119 and its analogues. For a comparison, the panels also include those in the Local Group and the other 9 quasar-intrinsic 2175 Å bump features. Among the 12 quasar-intrinsic cases, 3 objects (SDSS J0916+2921, SDSS J1020+6023 and SDSS J1349+3823) have redshifts of $z \approx 1.0$, and the others have higher redshifts ($z \approx 2.0$). The absorbers in the MW (blue circles) have a large absorption strength ($A_{\text{bump}} = 2.48 \pm 1.15 \mu\text{m}^{-1}$) and a constant peak position ($x_0 = 4.59 \pm 0.027 \mu\text{m}^{-1}$), and their widths vary over a wide range of $\gamma = 1.0 \pm 0.25 \mu\text{m}^{-1}$. The 2175 Å bump features in BAL quasars (pink squares; Zhang et al. 2015a) present broader bump features at smaller wave numbers ($\gamma = -1.3 \pm 0.39 \mu\text{m}^{-1}$, $x_0 = 4.45 \pm 0.05 \mu\text{m}^{-1}$), and their strengths are much weaker than those of the galactic bumps at 2175 Å on average, but they are similar to those of the LMC average (navy circles) and LMC2 Supershell (green circles) bumps. However, two 2175 Å bump features (brown triangles) reported in Pan et al. (2017) and Shi et al. (2020) are relatively prone to the MW sightlines that deviate most from the main distribution. The 2175 Å bump features in SDSS J1630+3119 and its analogues are significantly stronger than the previous detections in BAL quasars and SDSS J1349+3823. Their strengths are very close to the averaged strength of the galactic bumps. For the bump width, the bump in SDSS J1630+3119 has almost the same average width as those in the MW, and the bumps in the two analogues are slightly wider, but they are not beyond the range of those in the MW. For the peak position, only in one case (SDSS J1349+3823) does the bump peak at 2141 ± 4.5 Å: blueward of 2175 Å; the bumps in another two cases (SDSS J1020+6023 and SDSS J1630+3119) tend to peak at longer wavelengths, and their peak positions are 2227 ± 4.9 Å and 2217 ± 5.0 Å, respectively.

The bump peak position seems to vary in different studies. We note that this parameter depends on the accuracy of the redshift measurement, which is a weakness of BAL quasars. In their spectra, there is a lack of narrow emission lines and not enough complete BEL profiles to measure the precise redshift. The position deviation of the bumps in all BAL quasars is 0.08 μm$^{-1}$, and this value, corresponding to a redshift, is approximately 0.018. Although there are many disadvantages in the redshift
measurement of BAL quasars, the SDSS redshift accuracy is much higher than the above value. This suggests that the redshift uncertainty is not sufficient to cause such a large variation. Moreover, none of the 2175 Å bump features in this work and other quasar-intrinsic cases are beyond the range of the LMC average and LMC2 Supershell bumps (Gordon et al. 2003). Thus, the different parameter distributions of the 2175 Å bump features are probably intrinsic and may be attributed to variations in the grain-size and/or chemical composition of the dust (Draine 2003).

4.3. Survival environment of 2175 Å dust grains

The question for these FeLoBAL quasars with 2175 Å bump features is in which structure the 2175 Å dust grains survive: the BAL region, or something else. In the common sense, outflows may emerge from the outer region of the accretion disk or even from the innermost region of the torus, in which the gas clouds are dusty and relatively cold. These dusty clouds are lifted up and exposed to the central engine. Generally, the low-density part is rapidly heated and highly ionised, and the dust in this would be sublimated. The dust grains survive in dense regions and may even form in dense clouds embedded in outflows (Elvis et al. 2002). However, this does not mean that 2175 Å dust grains must exist in the outflows. Nearly two-thirds of the quasars with quasar-associated 2175 Å bump features are non-BAL quasars. When the sightline to the continuum source intersects the outflow clouds, the 2175 Å bump feature and BAL troughs would coexist in the spectrum, which originate from the dust grains and ionised gases carried by outflows, respectively. In addition, the torus may be the primary dust warehouse in the nuclear region of quasars. Large amounts of dust grains gather here, and the stacked structure provides a good shelter for them. The absorption features of the silicate grains at 9.7 µm and the polycyclic aromatic hydrocarbon (PAH) molecules at 3.4 µm are collected in the mid-IR spectra of many type II Seyferts. We suspect that this dust is widespread; it also certainly requires a proper orientation angle for 2175 Å bump features in quasars; otherwise, the dust grains either do not obscure the central engine or are completely obscure, and then the source is transformed into type II.

In Zhang et al. (2015a), we proposed another probable picture taken from the early evolution of quasars. The AGN activity destroys and clears the dust grains in the surrounding space of the nucleus, which are prepared through the persistent nucleosynthesis processes of carbon, and the formation of dust occurs throughout the evolution of stars and galaxies before the black hole is ignited. Outflow winds shield high-energy photons from the central engine and extend the survival period of dust in the nuclear region of quasars.

How these dust grains enter into (or are assembled) and survive in the nuclear region of quasars and the survival environment and the dust grain itself are interesting questions that deserve further observational study and investigation. For example, follow-up high-resolution spectroscopy at large telescopes will present abundant gas absorption lines in association with dust, which could be used to measure the chemical abundance and dust depletion, obtain the column density of hydrogen and the dust-to-gas ratio, and even detect neutral and molecular gases, so that we may finally understand the chemical and physical conditions in absorber systems.

Furthermore, although the origin of the 2175 Å bump features remains unclear, carbonaceous dust has long been proposed as a candidate. Laboratory and theoretical studies confirmed that materials such as graphite, aromatic hydrocarbon molecules, fullerene, hydrogenated amorphous carbon, and even carbon nanotubes could produce a broad extinction bump at approximately 2175 Å (e.g. Stecher & Donn 1965; Wright 1988; Mennella et al. 1998; Aannestad 1992; Joblin et al. 1992). The carrier of the 2175 Å bump features is probably one of them or a mixture of several materials. Follow-up near- and mid-IR spectroscopy will be expected to present the corresponding absorption and emission features of the materials.

5. Summary

We reported the discovery of 2175 Å bump features in the optical spectra of the FeLoBAL quasars SDSS J1630+3119 and its two analogues for the first time. They are typical FeLoBAL quasars, and their optical spectra present numerous absorption troughs from high- and low-ionisation species, such as the C IV, Al III, Mg II, and Fe II multiplets. After masking these broad absorption features, parametrised extinction curves were used to redden the quasar composite to fit the observed spectra. The UV to near-IR broad-band SEDs of these objects verify the efficiency of spectral continuum fitting. The fitting yields prominent extinction bumps at 2175 Å, which are also confirmed again through several independent validation tests. Compared to the bump features in the LMC and LMC2 Supershell, the detections in this work are much closer to those in the MW. Our 2175 Å bump features are detected in the quasar rest frame. Taking the fractions of high and (iron) low-ionisation BALs into account, we consider that they, as well as those in the BAL quasars of Zhang et al. (2015a), are intrinsic to quasars rather than to the quasar host and intervening galaxies. These 2175 Å bump features may be related to the bumps seen in the Local Group and may be a counterpart of the 2175 Å bump features under different conditions in the quasar environment. Traversing the various structures of AGNs, we suggest that only a dusty torus can provide a suitable context for the survival of 2175 Å bump carriers. The dust grains perhaps survive in the dusty torus, and the shielding effect of outflow clouds allows the MW-like dust to be assembled or to extend the survival period in the nuclear region. To further investigate the intrinsic nature of the FeLoBAL-2175 Å bump feature, we need follow-up monitoring of absorption lines and mid-IR spectroscopy to study PAH emissions.

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Appendix A: SDSS J1020+6023

SDSS J1020+6023 was selected as a candidate of the intrinsic 2175Å bump feature based on the SDSS spectrum. Its broad extinction bump was confirmed by UV-to-IR broad-band SED and again by follow-up Keck spectroscopy.

SDSS J1020+6023 was imaged at five photometric bands on March 3, 2000, and targeted as a quasar candidate based on its location in the uv′g′r′i′ colour cube (Richards et al. 2004) for spectroscopic observations by the SDSS survey. Two observations were taken on March 5, 2002, and the spectrum we used was extracted from the SDSS DR7. This object did not vary significantly between the SDSS photometric and spectroscopic observations because its spectrum is in excellent agreement with photometric data. SDSS J1020+6023 is also a typical FeLoBAL quasar, classified in several existing BAL catalogues (Trump 2006; Gibson et al. 2009; Zhang et al. 2010; Shen & Richards 2011). The blue-band spectrum shows overlapping iron absorption lines (2300-2415Å, 2550-2630Å and 2700-2770Å) and Mg II broad absorption troughs. Except for iron emission and absorption multiplets and narrow absorption lines of Mg I and Ca II H/K, no other emission or absorption lines can be identified from the spectrum.
extinction. After correcting for galactic reddening using the dust map of Schlegel et al. (1998) and the Fitzpatrick (1999) reddening curve, we transformed the magnitudes and optical spectrum into a rest frame with quasar redshift. The same fitting process was used to explore the 2175 Å bump feature in SDSS J1020+6023. The best-fitting model yields $c_1 = -1.82 \pm 0.02$, $c_2 = 0.54 \pm 0.01$, $c_3 = 0.95 \pm 0.08$, $x_0 = 4.49 \pm 0.01 \mu m^{-1}$, and $\gamma = 1.20 \pm 0.03 \mu m^{-1}$. The strength of the 2175 Å bump feature is $A_{\text{bump}} = 1.24 \pm 0.11 \mu m^{-1}$.

This model is consistent with multi-band photometry with effective wavelengths smaller than the H band. The GALEX NUV magnitude is high, consistent with the reddened quasar composite spectrum. The $u'$ band is affected by the expected strong Al III absorption and has a lower radiation flux than the model spectrum. However, the near-IR photometric fluxes ($H$, $K_s$, and $W1$) of SDSS J1020+6023 are below the reddened spectral template. The reason is that the near-IR composite was constructed from observations of 27 luminous, low-redshift, UV-excess-selected quasars, and it was excellently described by the combined model with a single power-law slope and a black body from Hα to 3.5 μm (Glikman et al. 2006). This result strongly suggests the presence of significant quantities of hot dust in the near-IR composite. We added a dust emission component with a temperature of $\sim 1700$ K to near-IR magnitudes of SDSS J1020+6023, and the correctional fluxes (open squares) are highly consistent with the near-IR composite. The typical dust sublimation temperature is $\sim 1500 - 2000$ K (e.g. Tuthill et al. 2001; Monnier & Millan-Gabet 2002), and the additional hot dust in this object is probably the inner edge of the torus. The combination spectrum together with the IR through optical-to-UV photometric data, the best-fitting model, and the observed and modelled extinction curves are presented in Fig. A.1.

Independent tests were employed to confirm the authenticity of the detection. As shown in Fig. A.2 (a), the bump strength distribution of the control sample peaks at $A_{\text{bump}} = 0.13 \mu m^{-1}$ with a variance of $\sigma = 0.16 \mu m^{-1}$, and the bump feature in SDSS J1020+6023 is detected at $> 5\sigma$. In Fig. A.2 (b), the scaled observed spectrum of the FeLoBAL quasar SDSS J112828.31+011337.9 is overplotted for comparison. This case has almost the same Mg II and Fe II absorption troughs as SDSS J1020+6023, and its spectrum agrees well with the corrected spectrum of SDSS J1020+6023. Both spectra follow the reddened quasar composite by the SMC extinction curve of $E(B-V) = 0.25$. 
Appendix B: SDSS J1349+3823

SDSS J1349+3823 is another analogue of SDSS J1630+3119. The archives provide multi-wavelength photometric magnitudes from the GALEX, SDSS, and 2MASS surveys and spectroscopic fluxes from the SDSS survey. Its 1D spectrum from 3800 Å to 9200 Å was accessed from the SDSS DR7. After correcting for galactic reddening using the dust map of Schlegel et al. (1998) and the Fitzpatrick (1999) reddening curve, we transformed the magnitudes and optical spectrum into the quasar rest frame. The photometric and spectroscopic data are shown using orange squares and black curves in Fig. B.1 (a). The optical spectrum simultaneously shows a remarkable intrinsic 2175 Å bump feature and broad absorption troughs of the Mg II and Fe II multiplets UV 1, UV 2, UV 3, and UV 62,63. The extremely low flux at the $u'$ band implies that there are possible strong absorption troughs of Al III and even Fe III UV $\lambda\lambda\lambda$ 1895.46, 1914.06, and 1926.30, and UV $\lambda\lambda\lambda$ 48 $\lambda\lambda\lambda$ 2062.21, 2068.90, and 2079.62.

The same fitting process was used to explore the 2175 Å bump feature in SDSS J1020+6023. The best-fitting model yields $c_1 = -0.73 \pm 0.02$, $c_2 = 0.37 \pm 0.01$, $c_3 = 2.02 \pm 0.13$, $x_0 = 4.67 \pm 0.01 \mu m^{-1}$, and $y = 1.18 \pm 0.08 \mu m^{-1}$, where the bump strength is $A_{\text{bump}} = 2.68 \pm 0.25 \mu m^{-1}$. This case, similar to SDSS J1020+6023, has a low redshift of $\sim 1.0$; therefore, the SDSS spectrum cannot cover the region shortward of the 2175 Å bump feature. However, the GALEX NUV and 2MASS near-IR photometric data fall right on the modelled spectrum, which suggests that the shape and bump of the extinction curve are well defined. In Fig. B.1 (a) and (b), the best-fitting model and observed and modelled extinction curves are presented. We also performed tests as for J1630+3119 to confirm the bump detection in SDSS J1349+3823. As shown in Fig. B.2, the bump feature in SDSS J1349+3823 is detected at $\gg 5\sigma$. The corrected spectrum of SDSS J1349+3823 agrees with the observed spectrum of the FeLoBAL quasar SDSS J084044.41+363327.8 and their spectral continua follow the reddened quasar composite by the SMC extinction curve of $E(B-V) = 0.17$. 

Fig. B.1. Same as Fig. 3, but for SDSS J1349+3823.

Fig. B.2. Same as Fig. 4 and Fig. 5, but for SDSS J1349+3823. The continuum slopes of the corrected spectrum of SDSS J1349+3823 and the observed spectrum of SDSS J084044.41+363327.8 follow the reddened quasar composite (red curve) by the SMC extinction curve of $E(B-V) = 0.17$. 

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