Overview of CNM LGAD results: Boron Si-on-Si and epitaxial wafers

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ABSTRACT: Low Gain Avalanche Detectors (LGADs) are n-on-p silicon sensors with an extra p-layer below the collection electrode which provides signal amplification. When the primary electrons reach the amplification region new electron-hole pairs are created that enhance the generated signal. The moderate gain of these sensors, together with the relatively thin active region, provide precise time information for minimum ionizing particles. To mitigate the effect of pile-up at the HL-LHC the ATLAS and CMS experiments have chosen the LGAD technology for the High Granularity Timing Detector (HGTD) and for the End-Cap Timing Layer (ETL), respectively. A full characterization of recent productions of LGAD sensors fabricated at CNM has been carried out before and after neutron irradiation up to $2.5 \times 10^{15} \text{n}_{eq}/\text{cm}^2$. Boron-doped sensors produced in epitaxial and Si-on-Si wafers have been studied. The results include their electrically characterization (IV and bias voltage stability) and performance studies (charge and time resolution) for single pad devices with a Sr-90 radioactive source set-up. The behaviour of the Inter-Pad region for irradiated 2 $\times$ 2 LGAD arrays, using the Transient Current Technique (TCT), is shown. The results indicate that the Si-on-Si devices with higher resistivity perform better than the epitaxial ones.

KEYWORDS: LGAD, Silicon sensors, Timing detector
1 Introduction

In the past years, several doping materials (Boron\textsuperscript{1}, Boron plus Carbon\textsuperscript{1} and Gallium doping\textsuperscript{2}) for LGAD sensors were explored by the CNM Barcelona\textsuperscript{[1]}. Their charge collection and time resolution as a function of the bias voltage at different temperatures were studied for both neutron and proton irradiation\textsuperscript{[2]}. Boron and Boron plus Carbon sensors perform similarly in terms of time resolution before and after irradiation (below 50 ps at higher voltages) but devices with Carbon collect more charge at the same bias voltage. Gallium presents 20\% less gain with respect to Boron doping and this research line has not been followed due to the poorer radiation hardness and timing performances. Carbon infusion seems to help more in diminishing the effect of gain reduction after irradiation. However, in this first run with Carbon the implantation process was not optimal and its benefits are not clear compared to later productions from other manufacturers.

2 CNM LGAD productions

More recently, CNM has produced two other LGAD runs with Boron doping for which the electrical characterization, the charge collection and timing performance are shown in this paper before and after neutron irradiation\textsuperscript{[3]}. The first run on Si-on-Si wafers\textsuperscript{[4]} with higher doping dose (1.8 ×

\textsuperscript{1}Production run 10478
\textsuperscript{2}Production run 10924
\textsuperscript{3}Irradiation for all the sensors has been performed with 1 MeV neutrons in the TRIGA reactor in Ljubljana.
\textsuperscript{4}Production run 12916, named AIDA2020.
10^{13} \text{ atoms/cm}^2\) than the first Boron run mentioned in Section 1. The second run is on low-resistivity epitaxial wafers\(^5\) with even a higher dose than the Si-on-Si devices \((2 \times 10^{13} \text{ atoms/cm}^2).\) The aim to investigate the usage of epitaxial wafers is their lower cost and the easier availability with respect to the Si-on-Si ones.

3 Electrical characterization

Si-on-Si sensors irradiated up to \(2.5 \times 10^{15} \text{n}_{eq}/\text{cm}^2\) and low-resistivity epitaxial sensors irradiated up to \(10^{16} \text{n}_{eq}/\text{cm}^2\) were electrically characterized in the CNM laboratory on a probe station at \(-30^\circ\text{C}\) in a dry environment. The leakage current was studied as a function of the bias voltage, in order to determine the breakdown voltage, \(V_{bd}\). In particular, gain layer and full depletion voltages, \(V_{gl}\) and \(V_{fd}\), are determined by C-V curves [3]. The breakdown voltage measured after dicing is similar to the one measured at a wafer level [3] despite the \(10^\circ\text{C}\) difference. Sensors were biased with positive voltage on the pad and guard-ring, while the back of the sensor (ohmic side) was grounded. The plots in Figure 1 show the total current \(I_{tot}=I_{PAD}+I_{GR}\) for the sensors from Si-on-Si (left) and low-resistivity epitaxial (right) wafers before and after irradiation. In particular, for Si-on-Si sensors \(V_{gl} \sim 38 \text{ V}, V_{fd} \sim 42 \text{ V}\) and \(V_{bd} \sim 50 \text{ V}\), and for epitaxial ones \(V_{gl} \sim 30 \text{ V}, V_{fd} \sim 35 \text{ V}\) and \(V_{bd} \sim 400 \text{ V}\) before irradiation. As expected, the bulk current and the breakdown voltage increase with the increasing of fluence for both productions. The breakdown voltage increases because of the lower quantity of dopant in the doping layer with respect to the less irradiated and unirradiated sensors [4].

![Figure 1. IV curves on probe station for Si-on-Si (left) and low-resistivity epitaxial (right) wafer sensors before and after irradiation.](image)

4 Stability measurements

An important parameter of the LGAD sensors is the voltage at which the rate of self induced triggers becomes large enough to prevent the operation of the devices, since this would generate a detector occupancy beyond the capability of the readout system. In the ATLAS High Granularity Timing Detector the rate of self triggers is limited to 1 kHz. The rate of self triggers for both Si-on-Si (left)
and low-resistivity epitaxial (right) sensors is shown in Figure 2. The unirradiated sensor from the Si-on-Si wafer presents a high rate of self triggers which severely limits the operational range of the devices before irradiation, since, as stated in the previous section, the full depletion is of the order of 42 V. However, after irradiation the operational range is extended. On the other hand, for the devices from the epitaxial wafers the voltage for self triggers is much larger, and it is possible to operate the sensors before and after irradiation. Epitaxial sensors needs more voltage at the same fluence, for $10^{14} \text{n}_{eq}/\text{cm}^2$ auto-triggers later than for Si-on-Si devices.

![Figure 2. Rate of self triggers measured at $-30 \degree \text{C}$ for Si-on-Si (left) and low-resistivity epitaxial (right) wafer sensors.](image)

5 Charge collection and time resolution

To study the performance of the sensors, coincidence triggers are acquired between the devices under test (DUTs) and a reference LGAD sensor with a Sr-90 radioactive source at $-30 \degree \text{C}$ in a dry environment. The setup is the same shown in [2] and the data are analyzed offline with the LGADUtils framework [5] developed at IFAE from which the charge collection and time resolution are calculated. The charge is computed for each event as the integral of the signal area. For each bias voltage point the charge distribution is fitted with a Landau-Gauss convoluted function and the collected charge is defined as the most probable value of the fit function. For the time resolution, the Constant Fraction Discriminator (CFD) method is used, taking into account the contribution of the reference sensor, which is $\sigma_{REF} = 35.7 \text{ ps}$, calculated from previous calibration measurements [6]. The $\sigma_{DUT}$ is obtained and presented as a function of the bias voltage for different neutron fluences.

5.1 Si-on-Si wafer

The collected charge and time resolution are shown in Figure 3. Sensors irradiated to the highest fluence reach the target\(^6\) of 4 fC for bias voltages higher than 680 V. All the irradiated sensors achieve time resolutions below 40 ps. Regarding the unirradiated sensor, this one shows a poor performance in time resolution due to the high rate of self-triggers already at 40 V, see Figure 2 (left). Comparing the time resolution value between the unirradiated sensor with the $10^{14} \text{n}_{eq}/\text{cm}^2$ irradiated sensor, for a fixed value of collected charge, the time resolution on the unirradiated device is poorer than the one for the irradiated sensor. This is due to the high rate of self trigger of the

\(^6\)The 4 fC low limit is given by the HGTD sensor and electronics requirements.
unirradiated sensor, which is, for a bias voltage value of $40 \, \text{V}$ and a charge collected of $\sim 19 \, \text{fC}$, around $7 \, \text{kHz}$. For the irradiated device, at $100 \, \text{V}$ bias voltage with a collected charge of $\sim 21 \, \text{fC}$, the rate of self triggers is below $1 \, \text{kHz}$, not affecting the timing performance of this device.

**Figure 3.** Collected charge (left) and time resolution (right) for Si-on-Si wafer sensors before and after irradiation.

### 5.2 Low-resistivity epitaxial wafer

The collected charge and time resolution are shown in Figure 4. The sensor irradiated at the highest fluence does not show any gain up to $720 \, \text{V}$, but the high leakage current (see Figure 1 (right)) prevented safely reaching higher voltages. The sensor irradiated at $10^{15} \, \text{neq/cm}^2$ reaches a collected charge of $4 \, \text{fC}$ for bias voltages higher than $700 \, \text{V}$. Unirradiated and irradiated sensors feature a time resolution around $40 \, \text{ps}$.

**Figure 4.** Collected charge (left) and time resolution (right) for low-resistivity epitaxial wafer sensors before and after irradiation.

### 6 Transient Current Technique

#### 6.1 Inter-pad distance for Si-on-Si sensors

The inter-pad region of $2 \times 2 \, \text{LGAD pads of } 1.3 \times 1.3 \, \text{mm}^2$ from the Si-on-Si wafer has been studied using the Transient Current Technique (TCT) [7]. The nominal gap before irradiation, that is the gap at the sensor design level, is $57 \, \mu\text{m}$. Four irradiated sensors from $10^{14}$ to $2.5 \times 10^{15} \, \text{neq/cm}^2$ have been measured at $-20 \, ^\circ\text{C}$ in a dry environment. An infra-red (IR) laser illuminated the sensor
through an opening in the metal layer on the backside of the device. A 0.5 × 0.5 mm² inter-pad area between two pads has been scanned with 50 µm steps in both X-Y directions collecting an average of 1000 waveforms for each laser position. Data have been analyzed offline and a 2D map has been built for each voltage point. A (reverse) error function is used to fit the 2D map projection distribution for each of the two pad of the DUT, as shown in Figure 5 (left). The inter-pad distance is estimated as the difference of the X positions at 50% of the height of the fit function for each of the two pads: \[ \text{IP}_{\text{gap}} = |x_{50\% \text{PA}D_{A}} - x_{50\% \text{PA}D_{B}}| \]. The IP\(_{\text{gap}}\) is shown in Figure 5 (right) as a function of bias voltage. At low fluences (blue line), part of the carriers generated underneath the gain layer drift to the Junction Termination Extension (JTE) and don’t produce charge multiplication, the effect is reduced as the bias voltage is increased. However, at higher fluence some gain is achieved by carries drifting at the JTE-p bulk interface due to field focusing and this results in smaller inter-pad distances than nominal value measured at higher voltages (black line). This behavior was also observed on LGADs from other vendors [8].

![Figure 5](image.png)

**Figure 5.** Projections and error function fits of the 2D map for each of the two pads (left) and inter-pad distances (right) measured with TCT for several fluences of Si-on-Si 2 × 2 LGAD arrays. The difference in collected charge of the two pads in the left figure is related to the non-perpendicular incidence of the infrared light on the entire sensor surface.

### 6.2 Gain and Time resolution for epitaxial sensors

Gain measurements have been performed using the TCT at −20 °C with an IR laser (Figure 6 left). Figure 6 (right) also shows the time resolution obtained from the time difference of a laser pulse and the same pulse delayed by 50 ns. These preliminary results are in agreement within the errors (not reported in the plot) with results measured with a Sr-90 radioactive source shown in Figure 4, considering the use of a different CFD fraction (usually > 50% for TCT).

### 7 Conclusion

Measurements of Si-on-Si devices show that the target charge of 4 fC can be achieved up to \(2.5 \times 10^{15} \text{n}_{eq}/\text{cm}^2\) at 700 V, while epitaxial devices require a higher voltage of 750 V for fluences up to \(10^{15} \text{n}_{eq}/\text{cm}^2\). The better performances of Si-on-Si could be explained by the difference of the wafer resistivity, since the low-resistivity (100 Ω cm) of the epitaxial wafers requires a higher bias to be operated. On the other hand, the Si-on-Si devices show a marginal performance before
irradiation due to the restricted operational voltage range. TCT inter-pad gap measurements show that after sufficient irradiation the effective gaps are lower than the nominal values. Preliminary results of time resolution using the TCT are in agreement with the ones obtained with the Sr-90 set-up for epitaxial sensors.

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