Spectral evolution of *Fermi*/*GBM* short gamma-ray bursts

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Accepted 2010 November 1. Received 2010 November 1; in original form 2010 August 27

**ABSTRACT**

We study the spectral evolution of 13 short-duration gamma-ray bursts (GRBs) detected by the Gamma Burst Monitor onboard *Fermi*. We study spectra resolved in time at the level of 2–512 ms in the 8 keV–35 MeV energy range. We find a strong correlation between the observed peak energy $E_{\text{peak}}$ and the flux $P$ within individual short GRBs. The slope of the $E_{\text{peak}} \propto P^\gamma$ correlation for individual bursts ranges between $\sim 0.4$ and $\sim 1$. There is no correlation between the low-energy spectral index and the peak energy or the flux. Our results show that in our 13 short GRBs, $E_{\text{peak}}$ evolves in time tracking the flux. This behaviour is similar to what found in the population of long GRBs and it is in agreement with the evidence that long GRBs and (the still few) short GRBs with measured redshifts follow the same rest-frame $E_{\text{peak}}-L_{\text{iso}}$ correlation. Its origin is most likely to be found in the radiative mechanism that has to be the same in both classes of GRBs.

**Key words:** radiation mechanisms: non-thermal – gamma-ray bursts: general.

1 INTRODUCTION

Short gamma-ray bursts (GRBs) have been a challenge since the finding of their spectral diversity with respect to the class of long GRBs (e.g. see Nakar 2007; Lee & Ramirez-Ruiz 2007 for recent reviews). Short GRBs have optical and X-ray afterglows, and in few cases, they also show X-ray flares, similar to those of long bursts but scaled by their fluence (e.g. Gehrels et al. 2008; Nousek, Fruchter & Pe’er 2009). However, the class of short GRBs is somewhat heterogeneous for what concerns the prompt emission properties (e.g. there are short GRBs followed by a faint extended emission – Norris & Bonnell 2006; Donaghy et al. 2006; Norris & Gehrels 2008) or the host galaxy properties (short GRBs are found in almost all galaxy types – e.g. Berger 2009). Recently, this picture has been also complicated by the detection of long GRBs at very high redshifts, like GRB 080913 at $z \sim 6.7$ (Greiner et al. 2009) and GRB 090423 at $z \sim 8.2$ (Salvaterra et al. 2009; Tanvir et al. 2009), which have an intrinsic duration of less than 2 s. Classification schemes of short GRBs that try to merge all this evidence have been proposed (Zhang 2007).

For what concerns the prompt emission, it was discovered through the BATSE sample that short GRBs are spectrally harder than long GRBs (Kouveliotou et al. 1993). A detailed analysis of the time-integrated spectra of short GRBs (Paciesas et al. 2001; Ghirlanda, Ghisellini & Celotti 2004; Ghirlanda et al. 2009) has also shown that this spectral difference is due to a harder low-energy spectral component in short GRBs (but see Nava et al. 2010). These studies also revealed that, on average, in the first 2 s of emission of long GRBs, they have similar spectral properties to that of short bursts (as also found from the comparison of the variability patterns – Nakar & Piran 2002), thus suggesting the presence of a common emission mechanism that operates in these two kinds of sources.

Considered individually, long GRBs show a strong spectral evolution with two possible behaviours: the peak energy $E_{\text{peak}}$ of the $\nu F_{\nu}$ spectrum decays in time (‘hard-to-soft’ evolution) or follows the variation in the flux (‘tracking’ evolution; e.g. Band et al. 1993; Ford, Band & Matteson 1995). This emerged from the time-resolved spectral analysis of long GRBs observed by the BATSE. Recently, Lu, Hou & Liang (2010) reported that the hard-to-soft spectral evolution is dominant in long bursts (at least initially, for two-thirds of their sample of 22 single-pulse BATSE GRBs). To date, however, no detailed study of the possible spectral evolution of short GRBs was performed.

Another recent issue concerning the prompt emission of GRBs is the nature of the spectral-energy correlations discovered for the sample of long GRBs with measured redshifts. Amati et al. (2002) found that the rest-frame peak energy correlates with the isotropic energy (the $E_{\text{peak}}-E_{\text{iso}}$ correlation), while Yonetoku et al. (2004) showed that a similar correlation exists between $E_{\text{peak}}$ and the isotropic luminosity (the $E_{\text{peak}}-L_{\text{iso}}$ correlation). The open debate is if these correlations have a physical foundation (Ghirlanda, Ghisellini & Frontera 2005; Bosnjak et al. 2008; Ghirlanda et al. 2008; Nava et al. 2008; Amati, Frontera & Guidorzi 2009; Krimm et al. 2009) or if they are the result of selection effects (Band & Preece 2005; Nakar & Piran 2005; Butler et al. 2007; Butler, Kocevski & Bloom 2009; Shahmoradi & Nemiroff 2010). Both the $E_{\text{peak}}-E_{\text{iso}}$ and $E_{\text{peak}}-L_{\text{iso}}$ correlations are defined by considering the time-integrated...
spectral properties of long GRBs with measured redshifts. Recently, Ghirlanda et al. (2009) showed that short GRBs with measured redshifts also follow the same $E_{\text{peak}}-L_{\text{iso}}$ correlation defined by long bursts, but not the corresponding $E_{\text{peak}}-E_{\text{iso}}$ correlation.

One of the most stringent results supporting a physical origin of the $E_{\text{peak}}-L_{\text{iso}}$ correlation is that, in long GRBs with measured redshifts detected by Fermi, there is a strong $E_{\text{peak}}-L_{\text{iso}}$ correlation within individual bursts (Ghirlanda, Nava & Ghisellini 2010), which cannot be due to selection effects but must have a physical origin. A similar result was reached by Firmani et al. (2009) based on the time-resolved spectral analysis of long GRBs detected by Swift. Several interpretations for the $E_{\text{peak}}-L_{\text{iso}}$ (or $E_{\text{peak}}-E_{\text{iso}}$) correlations have been proposed (e.g. Yamazaki, Ioka & Nakamura 2004; Eichler & Levinson 2005; Lamb, Donaghy & Graziani 2005; Levinson & Eichler 2005; Rees & Meszaros 2005; Tomo, Yamazaki & Nakamura 2005; Barbiellini et al. 2006; Ryde et al. 2006; Thompson 2006; Giannios & Spruit 2007; Thompson, Meszaros & Rees 2007; Guida, Bernardini & Bianco 2008; Panaitescu 2009), but there is no general consensus about a prevalent idea.

Therefore, the two topics we aim to address in this work are: (1) whether and how the spectra of short GRBs evolve in time; and (2) if also in short GRBs there is a spectral energy correlation between the peak energy and flux similar to what found in long events.

Time-resolved spectral studies of short GRBs are hampered by their lower fluence and duration with respect to long GRBs. Moreover, time-resolved spectroscopy requires an instrument with the largest energy coverage and good spectral resolution in the keV–MeV energy range in order to constrain the spectral parameters of individual (time-resolved) spectra and to follow their temporal variation. Swift detected a large number of short GRBs, but it is not suited to this aim because of the narrow energy range of the BAT instrument (15–150 keV). The Gamma Burst Monitor (GBM; 8 keV–40 MeV) onboard the Fermi satellite, instead, represents a unique opportunity to study, in details for the first time, how the spectrum in short GRBs evolves with time.

In Section 2, we present the sample of short Fermi GRBs selected for the time-resolved spectral analysis. The details of the spectral analysis are described in Section 3 and the results in Section 4. We draw our conclusions in Section 5.

2 THE SAMPLE

We consider the 237 GRBs detected by the Fermi GBM up to 2010 May and whose detection has been published through GCN communications. In this sample, there are 37 short GRBs with observed duration (as reported in the GCNs) $\leq 2$ s. For the time-resolved spectral analysis, we consider the short GRBs with the largest fluence and peak flux. This ensures to have, for each GRB, a set of time-resolved spectra with enough signal to constrain the spectral parameters. We select the 14 short GRBs with a fluence larger than $8 \times 10^{-7}$ erg cm$^{-2}$ (integrated in the 8 keV–1 MeV energy range) and a peak photon flux $F_{\text{peak}} \geq 11$ photon cm$^{-2}$ s$^{-1}$.

We anticipate that for one burst (GRB 100223), the detector response files are not present in the archive, so that its spectral analysis was not possible. GRB 081113 has a time-integrated spectrum, which is consistent with a single power law (von Kienlin et al. 2008a), and similarly its time-resolved spectra are all consistent with a single power law. All the other GRBs of our sample have time-integrated spectra (as reported in the GCN) with a peak in $vF_{v}$.

We have included GRB 090308 and GRB 081107 in the sample, although their duration ($T_{90}$) is somewhat longer than 2 s (i.e. 2.2 and 2.11 s, respectively), since they have a fluence and peak flux larger than our threshold and considering that the division of short and long GRBs at 2 s is not sharp, since both populations have a Gaussian $T_{90}$ distribution extending below and above this time cut-off. These two bursts are reported separately in Table 1. Therefore, the sample of analysed short GRBs is composed of 14 events.

In Table 1, we report the names, duration, fluence and peak flux of the selected short GRBs (with the corresponding reference – column 4). The redshift is not measured for all the bursts in our sample, except for GRB 090510 at $z = 0.903$ (Rau et al. 2009).

3 SPECTRAL ANALYSIS

The GBM (Meegan, Lichti & Bhat 2009) comprises 12 thallium sodium iodide [NaI(Tl)] and two bismuth germanate (BGO) scintillation detectors, which cover the energy ranges $\sim 8$ keV–1 MeV and $\sim 300$ keV–40 MeV, respectively. The GBM acquires three types of data suited for spectroscopic studies (Meegan et al. 2009): the ‘CTIME’ data consisting of a sequence of spectra (binned in eight energy channels) with a time-resolution between 0.256 and 1.024 s, the ‘CSPEC’ data containing a sequence of 128 energy channel spectra binned in time with a variable resolution between 1.024 and 4.096 s, and the ‘TTE’ event data files containing individual photons with time and energy tags.

To the aim of studying short GRBs, TTE data are ideal, because they allow to choose the temporal resolution according to the burst duration and flux. Due to the limited buffer size, TTE data only cover the interval between $\sim 30$ s before and $\sim 300$ s after the burst trigger time. This time-interval fully encompasses the duration of short GRBs and allows to fit the background spectrum by selecting time-intervals before and after the burst.

For the time-resolved spectral analysis, we used the recently released software RMFIT1 (v33pr7). In order to model the background spectrum for the time-resolved spectroscopic analysis, we selected two time-intervals before and after the burst. The sequence of background spectra in the two selected intervals was fitted with a first-order polynomial to account for the possible time-variation of the background spectrum. The background spectrum was then extrapolated to the time-intervals selected for the time-resolved spectroscopy of individual GRBs.

For each burst, we jointly fitted the spectra from the NaI detectors, which had the largest illumination by the GRB. The NaI detectors were selected for having an angle of position of the GRB with respect to the detector normal lower than $\sim 80^\circ$. Among the two BGO detectors, we selected that with the smallest angle to the GRB. Four NaI detectors and one BGO detector were analysed for GRB 100206, 090902, 081213, 090108, 090305, 090308, 081216, 090228 and 090227B; three NaI detectors and one BGO detector for GRB 090206 and 090328; and three NaI detectors for GRB 081107. For GRB 090510, we used five NaI and both BGO data. The inclusion of the BGO data extends the spectral coverage of the NaI detectors from 1 MeV to $\sim 35$ MeV. This energy extension represents an unprecedented opportunity for the spectral analysis of short Fermi GRBs whose time-integrated spectrum has typically a peak energy of the order of 1 MeV (e.g. Nava et al. 2010).

For each burst we performed a time-resolved spectral analysis by opportunistically changing the time-resolution (starting with 64 ms time-resolution of the TTE data) and we checked that each time-resolved spectrum had well-constrained spectral parameters (i.e. the normalization, $E_{\text{peak}}$ and $\alpha$ are required to have relative errors

1 http://fermi.gsfc.nasa.gov/ssc/data/analysis/user/
Table 1. Short GRB sample. The last three digits in parentheses are the fractional trigger number assigned to the GRB by the GBM archive. The fluence $F_{-6}$ is in units of $10^{-6}$ erg cm$^{-2}$. References for the duration ($T_{90}$), fluence ($F_{-6}$) and peak flux ($P_{\text{peak}}$); (1) Guiriec et al. (2009a); (2) Guiriec et al. (2009b); (3) von Kienlin et al. (2009); (4) McBreen et al. (2008); (5) Wilson-Hodge et al. (2009); (6) Bissaldi et al. (2009b); (7) Goldstein et al (2009a); (8) Bissaldi et al. (2008); (9) von Kienlin et al. (2008b); (10) Bissaldi et al. (2009a); (11) Goldstein et al (2009b); (12) von Kienlin et al. (2010); (13) Goldstein et al. (2009c); and (14) von Kienlin et al. (2008a) – column 6 lists the range of time-resolution in milliseconds used to perform the spectral analysis. In the case of GRB 0909228, the time-resolution is 2 ms for the first peak and 32 ms for the (fainter) second peak of its light curve.

| GRB         | $T_{90}$ (s) | $F_{-6}$ (erg cm$^{-2}$) | $P_{\text{peak}}$ (photon cm$^{-2}$ s$^{-1}$) | Reference | $T_{\text{res}}$ (ms) |
|-------------|-------------|-------------------------|-----------------------------------|-----------|-------------------|
| 090510(016)$^a$ | 1.0         | 30.4 ± 2.0              | 80                                | 1         | 16–64             |
| 090227(772)  | 0.9         | 8.7 ± 0.1               | 34.6 ± 0.3                        | 2         | 8–32              |
| 090228(204)  | 0.8         | 6.1 ± 0.09              | 133 ± 8                           | 3         | 2–32              |
| 081216(531)$^a$ | 0.96       | 3.6 ± 0.1               | 55 ± 3                            | 4         | 16–32             |
| 090305(052)  | 2.0         | 2.7 ± 0.2               | 11 ± 2                            | 5         | 64–0.448          |
| 090902(401)  | 1.2         | 2.11 ± 0.14             | 11.4 ± 1.3                        | 6         | 64–0.512          |
| 091008(020)  | 0.9         | 1.28 ± 0.24             | 39.7 ± 3.9                        | 7         | 16–0.384          |
| 081223(419)  | 0.89        | 1.2 ± 0.10              | 22 ± 3                            | 8         | 64–0.704          |
| 081113(230)  | 0.5         | 1.07 ± 0.03             | 20 ± 1                            | 9         | ...               |
| 090206(620)  | 0.8         | 1.04 ± 0.06             | 19 ± 1                            | 10        | 64                |
| 090328(713)  | 0.32        | 0.96 ± 0.03             | 29.83 ± 2.38                      | 11        | 32–64             |
| 100206(563)  | 0.13        | 0.93 ± 0.04             | 31 ± 2                            | 12        | 32                |
| 090308(734)  | 2.11        | 3.46 ± 0.13             | 14.22 ± 0.91                      | 13        | 64                |
| 081107(321)  | 2.2         | 1.64 ± 0.28             | 11 ± 3                            | 14        | 64                |

$^a$GRBs with a precursor. Note: GRB 090308 and GRB 081107 are shown separately, since they have a slightly larger duration than the canonical 2 s dividing line between short and long GRBs, but still have a fluence and peak flux above our selection thresholds.

less than 100 per cent). For the brightest GRBs of our sample (i.e. GRB 090227B and 090228), we performed a time-resolved spectral analysis down to the 8- and 2-ms time-scale. In some cases, the end of the slow decline of the light curves did not ensure enough signal to obtain a good fit so that a coarser time-resolution was applied (e.g. a single bin of 0.5 s at the end of GRB 090902 was analysed).

In Table 1 (last column), we report the minimum and maximum time-resolution in milliseconds at which the time-resolved spectral analysis was performed.

The model adopted is a power law with an exponential cut-off whose free parameters are the low-energy spectral index $\alpha$, the peak energy $E_{\text{peak}}$ (i.e. the peak of the $\nu F_\nu$ spectrum) and the normalization. We also allowed for a variable normalization factor in order to fit together the data of the NaI and BGO detectors. For the time-resolved spectral analysis, we fixed the normalization factors to the values obtained from the fit of the time-integrated spectra of each burst. The values of these normalization factors are all consistent with 1, with a deviation of 20 per cent at most. The choice of this model is motivated by the fact that at high energies the response of the BGO detector rapidly decreases for increasing energy, so that it is hard, in single time-resolved spectra, to constrain the possible presence of a power-law component of, for example, the Band function (Band et al. 1993), which is instead typically fitted to the time-integrated spectra of long GRBs. We tried to fit the time-resolved spectra with the Band function, but in most cases, the high-energy spectral index $\beta$ was unconstrained. This is also confirmed by the recent analysis of the brightest three short GRBs of our sample (Guiriec et al. 2010) that shows that the fit with a Band function results in unconstrained $\beta$ in most of their time-resolved spectra. Therefore, for homogeneity and with the aim of comparing the spectral evolution trends in short GRBs, we adopted the same spectral model, that is, a cut-off power law for all the time-resolved spectra.

4 RESULTS

We searched for possible correlations between the best-fitting spectral parameters of the time-resolved spectra of our sample of short GRBs. We find that there is a strong correlation between the peak energy $E_{\text{peak}}$ and the flux $P$ within individual short GRBs. The flux $P$ is calculated by integrating the best-fitting model of each time-resolved spectrum in the 8 keV–35 MeV energy range. This correlation is shown in Fig. 1 for the whole sample of 13 GRBs. We could extract for them a total of 163 time-resolved spectra. Considering the 13 short GRBs individually, the $E_{\text{peak}}$–$P$ correlation of individual GRBs has a slope that ranges between 0.4 and 1.0. The $E_{\text{peak}}$–$P$ correlation is in the observer frame (only GRB 090510 has a known redshift). For comparison with the population of long GRBs, we show in Fig. 1 the $E_{\text{peak}}$–$L_{\text{iso}}$ (‘Yonetoku’) correlation of long GRBs with measured redshifts (adapted from Ghirlanda et al. 2010) assuming a redshift $z = 1$ and 0.5 (solid black and grey lines, respectively). Also shown are the boundaries representing the $3\sigma$ scatter of this correlation (dashed lines in Fig. 1). The short GRBs’ spectral evolution trends are consistent with this correlation (transformed in the observed frame) as already shown for the spectral evolution tracks of long GRBs (Ghirlanda et al. 2010). However, we note from Fig. 1 that the $E_{\text{peak}}$–$P$ correlation defined by short GRBs is similar in slope to that of the $E_{\text{peak}}$–$L_{\text{iso}}$ correlation transformed in the observer frame (solid grey and black lines in Fig. 1), but it has a higher normalization.

The most sampled events (i.e. those studied on the shortest integration time-scales – down to 4 ms in the case of the peak of GRB 090228) are the four GRBs with the largest fluence/peak flux (Table 1), shown separately in Fig. 2. It is noteworthy that two of them also show precursor activity, which is shown with different symbols (colours) in Fig. 2. GRBs are known to show, in different symbols (colours) in Fig. 2. GRBs are known to show, in different symbols (colours) in Fig. 2. GRBs are known to show, in different symbols (colours) in Fig. 2. GRBs are known to show, in different symbols (colours) in Fig. 2.
Figure 1. Correlation between the peak energy $E_{\text{peak}}$ and the flux of the 163 time-resolved spectra of the 13 short GRBs analysed in this work (Table 1). Different symbols/colours correspond to different bursts (as shown in the legend). The solid line (dashed lines) is the ‘Yonetoku’ relation (and its 3σ scatter) of long GRBs with measured redshifts (adapted from Ghirlanda et al. 2009) transformed in the observer frame assuming $z = 1$. The grey lines (solid and dashed) are for $z = 0.5$.

Figure 2. Correlation between $E_{\text{peak}}$ and the flux $P$ for the four GRBs with the highest fluence. The upper panel shows GRB 090510 and GRB 090227B. The precursor of GRB 090510 is shown with the green arrow symbols, since its spectrum is consistent with a single power law. The lower panel shows GRB 081216 (with a precursor, green symbols) and GRB 090228. Colour code (data points, solid and dashed lines) is the same as in Fig. 1.

preceding the main episode (i.e. precursors). Burlon et al. (2009) analysed a large sample of bright long BATSE GRBs with precursor emission and showed that the same spectral evolution is present in the precursors and in the main GRB events. They also did not find any relation between the spectral index and $E_{\text{peak}}$, similarly to what shown in Fig. 3. Here, we find that, although still based only on two cases, also in short GRBs the precursors have spectra, which are consistent with the general trend of the main event. This result, shown for short GRBs in this Letter for the first time, points towards a common origin of the precursor and the main emission.

Fig. 3 shows that there is no correlation between the low-energy spectral index $\alpha$ of the cut-off power-law model and the peak energy $E_{\text{peak}}$. Similarly, we do not find any correlation between $\alpha$ and the flux $P$. We instead find, as already shown by the analysis of time-integrated spectra of short GRBs (e.g. Ghirlanda et al. 2009), that they are harder than long ones and this makes a large fraction of them inconsistent with the ‘lines of death’ of synchrotron emission (with no cooling – vertical dotted line in Fig. 3) and makes all of...
5 CONCLUSIONS

The time-resolved spectra of individual short GRBs evolve in time and their $E_{\text{peak}}$ is strongly correlated with their flux. Furthermore, the found correlation is very similar to what already found in individual long GRBs. The GBM data for the brightest bursts allowed to analyse their spectra even at the 2–8 ms time-resolution for part of the duration of these GRBs. With respect to typical peak fluxes measured on ∼1 s time-bins, the peak fluxes reached by our short GRBs, measured on a fraction of second, are extreme, that is, of the order of several times $10^{-4}$ erg cm$^{-2}$ s$^{-1}$ (see Fig. 1).

We do not know if these flux levels are also reached in long GRBs, since a time-resolved analysis with this degree of accuracy has not yet been done (but time-resolved spectra with a coarser time-resolution for the brightest long bursts had much smaller peak fluxes e.g. Kaneko et al. 2006). The fact that the $E_{\text{peak}}$–$P$ correlation is similar suggests that long and short GRBs share the same emission mechanism for their prompt emission. Furthermore, the emission mechanism should not depend on the progenitor (e.g. fireball–funnel interactions), if long and short GRBs do have different progenitors.

For the brightest bursts, the dynamic range of the $E_{\text{peak}}$–$P$ correlation is very large, being more than two orders of magnitude in flux and more than one order of magnitude in $E_{\text{peak}}$. Even precursors obey the same $E_{\text{peak}}$–$P$ correlation, although only two short GRBs in our sample show a precursor.

The three brightest short bursts in our sample reach values of $E_{\text{peak}}$ significantly larger than those of long GRBs (e.g. Ghirlanda et al. 2010). On the other hand, they also reach significantly larger fluxes while remaining on the same $E_{\text{peak}}$–$P$ correlation defined by long bursts. For the same fluxes, they do have the same $E_{\text{peak}}$.

The $E_{\text{peak}}$–$P$ correlation, as well as the analogous Yonetoku relation for time-integrated spectra of different GRBs, has not yet received a convincing and broadly accepted explanation, as mentioned in the introduction. The fact that short bursts also obey it makes the search for a convincing interpretation even more compelling.

ACKNOWLEDGMENTS

We thank the referee for useful comments and suggestions. We acknowledge ASI (I/088/06/0) and a 2010 PRIN–INAF grant for financial support. DB is supported through DLR 50 OR 0405. This research has made use of the data obtained through the High Energy Astrophysics Science Archive Research Center Online Service provided by the NASA/Goddard Space Flight Center. DB and LN thank the Brera Observatory for the kind hospitality during the completion of this work.

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