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It’s not easy being green: The evolution of galaxy colour in the EAGLE simulation

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ABSTRACT

We examine the evolution of intrinsic $u$--$r$ colours of galaxies in the EAGLE cosmological hydrodynamical simulations, which has been shown to reproduce the observed redshift $z = 0$ colour-magnitude distribution well, with a focus on $z < 2$. The median $u$--$r$ of star-forming ('blue cloud') galaxies reddens by 1 magnitude from $z = 2$ to 0 at fixed stellar mass, as their specific star formation rates decrease with time. A red sequence starts to build-up around $z = 1$, due to the quenching of low-mass satellite galaxies at the faint end, and due to the quenching of more massive central galaxies by their active galactic nuclei (AGN) at the bright end. This leaves a dearth of intermediate-mass red sequence galaxies at $z = 1$, which is mostly filled in by $z = 0$. We quantify the time-scales of colour transition finding that most galaxies spend less than 2 Gyr in the ‘green valley’. We find the timescale of transition to be independent of quenching mechanism, i.e. whether a galaxy is a satellite or hosting an AGN. On examining the trajectories of galaxies in a colour-stellar mass diagram, we identify three characteristic tracks that galaxies follow (quiescently star-forming, quenching and rejuvenating galaxies) and quantify the fraction of galaxies that follow each track.

Key words: galaxies: colours, galaxies: evolution, galaxies: formation.

1 INTRODUCTION

A scatter plot of observed galaxies in optical colour versus broad-band magnitude (or stellar mass) reveals two relatively well-defined distinct populations, a ‘red sequence’ and a ‘blue cloud’, in a volume-limited sample. While a narrow red sequence was evident in early datasets (e.g. Sandage & Visvanathan 1978; Larson et al. 1980; Bower et al. 1991), this striking colour ‘bi-modality’ was perhaps first revealed most clearly by Strateva et al. (2001), who exploited the step-change in sample size offered by the Sloan Digital Sky Survey (SDSS, York et al. 2000), the statistics and general properties of these two sequences were subsequently characterised quantitatively by e.g. Baldré et al. (2004). The $u$--$r$ colour of a galaxy correlates strongly with its morphology (Hubble type); the GALAXY Zoo$^1$ citizen science project and the MEGAMORPH survey have put this correlation on a firm statistical footing (Willett et al. 2013; Häußler et al. 2013). Blue colours are typically due to light from massive hot, young stars. Broadly speaking, the more massive blue galaxies are star-forming discs and the fainter ones are irregulars. The red galaxies, in contrast, are elliptical or lenticular types, with the red colour reflecting an old stellar population. These galaxies are often referred to as ‘red-and-dead’ to stress that their star formation has mostly ceased (e.g. Brammer et al. 2009).

At higher redshifts, $z > 1$, the blue cloud is clearly dominant and its colour becomes increasingly blue with increasing $z$. Selection effects bias against the detection of red galaxies at higher $z$, but it is clear that the bright end of the red sequence is already in place by $z \sim 1$, albeit with a small shift in colour (e.g. Wolf et al. 2003; Bell et al. 2004)

It is difficult to establish how these sequences arise, or how individual galaxies evolve in colour space, using observations alone. This is because both star formation and galaxy destruction by mergers change the number density of galaxies of given mass across time. Faber et al. (2007) noted that the number density of blue galaxies remains approximately constant below redshift $z \sim 1$ whereas that of red galaxies increases markedly. This led them to propose
a model in which one or more mechanisms operate that decrease the star formation rate of blue galaxies, with such ‘quenching’ making galaxies redder until they join the red sequence. Bell et al. (2012) showed that there is significant scatter in the properties of quenched galaxies. One correlation that stood out in their sample, is that quenched galaxies usually exhibit a prominent bulge, by association suggesting a supermassive black hole.

A model in which an accreting supermassive black hole quenches star formation in its host galaxy appears very attractive. This is because black hole mass increases rapidly as a function of bulge mass (e.g. Haring & Rix 2004; McConnell & Ma 2013), hence such a model might explain why most massive galaxies are red - as observed. Unfortunately the evidence that star formation in galaxies hosting X-ray bright AGN is indeed suppressed appears inconclusive. Several studies have found no correlation between star formation rate and X-ray luminosity for an X-ray selected sample of AGN (e.g. Rosario et al. 2012; Harrison et al. 2012; Stanley et al. 2015). However a close to linear correlation that stood out in their sample, is that quenched galaxies usually exhibit a prominent bulge, by association suggesting a supermassive black hole.

Another well-documented process that quenches star formation in a galaxy is restriction of its supply of gas by either ram-pressure stripping of disc gas (Gunn & Gott 1972, e.g.) or removal of halo gas (e.g. strangulation Larson et al. 1980), as the galaxy traverses a region of higher gas pressure associated with a group or cluster. The quenching of star-formation turns these satellites red (e.g. Knobel et al. 2013). Originally suggested by Gunn & Gott (1972), the efficiency of these mechanisms has been investigated using simulations by many groups (e.g. Quilis et al. 2000; Roediger & Brüggen 2007), with more recently Bahé et al. (2013) pointing out that galaxies may be stripped before they become satellites, by the gas in the outskirts of massive systems. McCarthy et al. (2008) presented a theoretical framework that improves upon the simple analysis by Gunn & Gott (1972), and describes their simulation results well.

Observational confirmation that environmental quenching indeed operates is evidenced by the fact that red galaxies preferentially reside in regions of high galaxy number density (Dressler 1980), or equivalently that red galaxies are more strongly clustered than blue galaxies, even at fixed mass (e.g. Zehavi et al. 2005), and that the clustering amplitude of red galaxies depends little on mass (in contrast to that of blue galaxies, e.g. Coil et al. 2008). This is compatible with a model where red galaxies reside close to, or even inside, more massive and hence strongly clustered halos that cause the quenching. Trends between the environment and gas content of galaxies provide further evidence, with galaxies residing in clusters seen to be deficient in both HI and H2 gas relative to the field (e.g. Cortese et al. 2011; Boselli et al. 2014). Particularly convincing is the similarity of the trails of HI gas seen to be emanating from gas rich galaxies in clusters (e.g. Chung et al. 2007; Fumagalli et al. 2014) and of the ram-pressure stripped gas behind simulated galaxies that fall onto a cluster (e.g. Roediger & Brüggen 2008).

Even though observations suggest two empirical models of quenching (i.e. AGN and environmental), models of galaxy formation have struggled to reproduce simultaneously the detailed distribution of galaxies in the colour-magnitude diagram and the different clustering properties of red and blue galaxies. This is true of semi-analytical models, which use phenomenological prescriptions to describe the physical processes that lead to quenching (e.g. Font et al. 2008; Lacey et al. 2015); for example, Henriques et al. (2015) compare the Munich semi-analytical L-galaxies model to SDSS data. Although in many aspects this model reproduces the observations better than its predecessors, limitations remain. For example, L-galaxies’ u-r colours are considerably more bimodal than observed.

Hydrodynamical simulations can, in principle, model many physical processes self-consistently, but lack of numerical resolution and other limitations of the hydrodynamical integration may limit their realism. Fortunately, relatively small changes to the basic hydrodynamics scheme (e.g. Price 2008; Hopkins 2013) seem to resolve most numerical issues, such that the dominant uncertainties in hydrodynamical simulations become associated to the implementation of unresolved ‘subgrid’ processes rather than the details of the hydrodynamics scheme (Scannapieco et al. 2012; Schaller et al. 2015b).

The huge dynamic range required to simulate a cosmologically representative volume with the required resolution to follow the hierarchical build-up of galaxies, presents a major challenge to numerical simulations. Until recently, such simulations did not reproduce the galaxy stellar mass function well, let alone the detailed colours/clustering of galaxies. A red/blue bimodality appears in the zoomed-simulations of Cen (2014) even though these do not include AGN. However, the r−band luminosity function of these simulation contains many more massive galaxies than observed. Gabor & Davé (2012) include the effects of AGN using a heuristic prescription of heating gas, where cooling is simply switched off in halos deemed massive enough to host AGN. They illustrate how this process builds-up a red sequence below redshift z ~ 2; initially lower-mass satellites and more massive quenched centrals appear in heated halos, with a characteristic dip in the abundance of red galaxies of stellar mass M∗ ~ 10^9 M⊙ that is more prominent at higher z. While this simulation may provide valuable insight into the build up of the red sequence, the heuristic nature of the halo heating limits their practical applicability. For lower mass galaxies, Sales et al. (2015) show that the Illustris simulation (Vogelsberger et al. 2014) broadly reproduces the colours of satellites, which they attribute to the relatively large gas fractions of satellites at infall.

The Eagle reference model was calibrated to the z = 0.1 stellar mass function, black hole masses and sizes of galaxies and is currently the only hydrodynamical simulation that reproduces these observations. Eagle also reproduces many independent galaxy observations, such as the content and ionisation state of gas (Bahé et al. 2016; Lagos et al. 2015b), mass profiles (Schaller et al. 2015a) and evolution in stellar mass, star formation rate and size (Furlong et al. 2015b,a). The clustering of galaxies as a function of colour is investigated in a companion paper (Artale et
al., in prep.). Trayford et al. (2015) showed that EAGLE reproduces the $g-r - M_*$ colour magnitude (and the $g-r - M_\odot$) relation from the GAMA spectroscopic survey (Driver et al. 2011) very well. Including a model for dust-reddening computed using the SKIRT radiative transfer scheme (Baes et al. 2005; Camps & Baes 2015) improves the quantitative agreement further (Trayford et al. in prep.). With low-redshift ($z \sim 0.1$) galaxy colours in EAGLE appearing to be realistic, studying how they have arisen given the physical feedback model of the simulation may provide new insight. The evolution of EAGLE galaxy colours is also afforded credibility by the reasonable evolution of the EAGLE galaxy population in terms of the stellar mass function (Furlong et al. 2015b).

In section 2 we describe the EAGLE simulations used in this study, particularly the aspects of star formation, metal enrichment and feedback that are most relevant for setting intrinsic galaxy colours. We investigate in section 3 the evolution of the galaxy population across the colour-mass diagram and correlate colour changes with galaxies becoming satellites or hosting an AGN. In section 4 we expound these processes by analysing the behaviour of individual galaxies, using galaxy merger trees. Typical time-scales associated with colour transition are presented in section 4.2. We show that the colour evolution of most galaxies can be described well in terms of three generic tracks and quantify the fraction of galaxies that follow each path. Finally, our findings are summarised in section 5. Throughout this work we refer to dust-free, rest-frame colours as ‘intrinsic’ colours, and we take $Z_\odot = 0.0127$ for the metallicity of the Sun (Allende Prieto et al. 2001). Note that while the $Z_\odot$ value affects the normalisation of metallicities in solar units, colours are unaffected by the assumed $Z_\odot$ (see Trayford et al. 2015).

2 THE EAGLE SIMULATIONS

The EAGLE suite (Schaye et al. 2015; Crain et al. 2015) includes simulations performed in a range of periodic volumes and at various numerical resolutions to enable convergence testing. The simulations were performed with the GADGET-3 tree-SPH code (Springel 2005), but with changes to the SPH and time-stepping algorithm collectively referred to as ANARCHY (see appendix of Schaye et al. (2015) for details and Schaller et al. (2015b) for the relatively minor impact of these changes on the properties of simulated galaxies). We use the ΛCDM cosmological parameters advocated by Planck Collaboration et al. (2014), and initial conditions generated at $\z = 127$ (Jenkins 2013) using second order Lagrangian perturbation theory. We concentrate here on analysing the largest reference model (Ref-100). This is a cubic cosmological volume of 100 comoving Mpc (cMpc) on a side, with an initial gas particle mass of $m_g = 1.81 \times 10^5 M_\odot$. The simulation has a Plummer equivalent gravitational softening of $\epsilon_{\text{prop}} = 0.7$ proper kpc (pkpc) at redshift $z = 0$.

2.1 Subgrid model and galaxy identification

The EAGLE reference model implements subgrid modules for physical processes that occur below the resolution limit, corresponding approximately to the Jeans length of the warm ISM. The free parameters that enter the modules for feedback were calibrated using the redshift $z = 0.1$ galaxy stellar mass function, the $z = 0.1$ stellar mass-size relation, and the $z = 0$ stellar mass - black hole mass relation, see Crain et al. (2015) for motivation and details. We briefly summarise these subgrid modules here, paying particular attention to those aspects most crucial for this paper.

- Radiative cooling and photo-heating of gas by the evolving optically-thin UV/X-ray background of Haardt & Madau (2001) is implemented element-by-element following Wiersma et al. (2009a).

- Star formation is implemented as a pressure-law (Schaye & Dalla Vecchia 2008) so that simulated galaxies reproduce the observed $z = 0$ relation between gas and star formation surface density of Kennicutt (1998). Each gas particle is assigned a star formation rate, $\dot{m}_\star$, and gas particles are converted to star particles stochastically. The star formation rate is zero for particles below the metallicity-dependent threshold of Schaye (2004). We resample young stars using a probability proportional to the estimates of $\dot{m}_\star$, to reduce sampling noise when estimating the galaxy luminosities, as described in Trayford et al. (2015).

- Feedback from star formation is implemented by heating gas particles neighbouring newly-formed star particles as described by Dalla Vecchia & Schaye (2012). In this purely thermal implementation, a temperature boost $\Delta T_{\text{SF}}$ is defined and the heating of neighbouring gas particles is sampled stochastically given the energy available for feedback. A value of $\Delta T_{\text{SF}} = 10^{6.2}$K is chosen to be high enough to mitigate rapid cooling due to numerical effects, but low enough to avoid poor sampling of heating events around individual star particles (Dalla Vecchia & Schaye 2012).

- Seeding, merging, accretion and feedback from supermassive black holes is implemented as described in Schaye et al. (2015). Briefly, dark matter halos with virial mass $> 10^{10} h^{-1} M_\odot$ are seeded with a black hole of mass $10^5 h^{-1} M_\odot$. These can grow through Eddington-limited accretion of gas while accounting for the gas angular momentum as by Rosas-Guevara et al. (2015), and through mergers with other black holes, following Springel et al. (2005) and Booth & Schaye (2009). Feedback from accreting black holes is also modelled by heating surrounding gas using an implementation similar to that of stellar feedback. The temperature boost for AGN heating events is chosen to be $T_{\text{AGN}} = 10^{8.5}$K for all the simulations considered here.

Each star particle represents a ‘simple stellar population’ (SSP), characterised by an assumed stellar initial mass function (IMF, EAGLE adopts the Chabrier (2003) IMF over the mass range $[0.1, 100] M_\odot$), and assuming that stars have a metallicity inherited from the converted gas particle, with a single age corresponding to the time that the gas particle was converted to a star particle. We then use published stellar life-times, evolutionary tracks, and yields to compute the rate at which these stars evolve and lose mass, as well as the rate of core collapse and Type Ia supernova events as described in Wiersma et al. (2009b). The simulation tracks 11 elements (H, He, C, Ni, O, Ne, Mg, Si, S, Ca, Fe) as well as a ‘total metallicity’ (metal mass fraction) variable for each gas and star particle. The mass, age, and metallicity of the SSP are input parameters for the population synthesis model described below.
• Dark matter halos are identified using the ‘friends-of-friends’ algorithm (FOF), linking dark matter particles within 0.2 times the mean inter-particle separation into a single FOF halo. Other particles are assigned to the same halo (if any) as the nearest dark matter particle. We characterise the mass of the halo by its $M_{200,\text{crit}}$ value. This is the mass enclosed within a sphere of radius $R_{200,\text{crit}}$ centred on the location of the particle with minimum gravitational potential in the halo. This radius is chosen such that the mean density within this sphere is 200 times the critical density, given the assumed cosmology.

• Galaxies are identified with the SUBFIND algorithm (Springel et al. 2001; Dolag et al. 2009). SUBFIND identifies self-bound substructures within halos which we associate with galaxies. The ‘central’ galaxy is the galaxy closest to the centre of the parent FOF halo; this is nearly always also the most massive galaxy in that halo. The other galaxies in the same halo are its satellites. Particles in a halo not associated with a bound substructure (i.e. satellites) are assigned to the central galaxy. Central massive galaxies ($M_\star \geq 10^{11} \, M_\odot$, say) then have an extended halo of stars around them, usually referred to as intra-group or intracluster light. Determining the mass or indeed luminosity of such a large galaxy is ambiguous, both in simulations and in observations. For this reason we impose an aperture on the definition of a galaxy; we follow Schaye et al. (2015) and calculate masses and luminosities for every subhalo, excluding material that is outside a 30 pkpc spherical aperture centred on the subhalo potential minima as well as material that is not bound to that subhalo. The 30 pkpc aperture has been shown to mimic an observational Petrosian aperture, and reduces intra-cluster light in massive centrals while lower mass galaxies are unaffected (Schaye et al. 2015).

2.2 Galaxy colours

The stellar population properties (age, metallicity & assumed IMF) of an EAGLE galaxy are combined with the Bruzual & Charlot (2003) population synthesis model to construct an SED for each star particle. Summing spectra over all stars within the aperture described in section 2.1 and convolving with a filter response function yields broad-band colours, which we compute using the ugrizYJHK photometric system for optical and near infrared photometry (taken from Dai et al. 2010; Hewett et al. 2006). We express these absolute magnitudes in the AB-system, see Trayford et al. (2015) for more details.

It is well known that dust can alter the optical colour of a galaxy significantly, particularly for gas-rich discs seen edge-on. We describe a simple model for dust reddening in a previous study (Trayford et al. 2015), as well as a model that uses ray-tracing to account for the patchy nature of dust clouds enshrouding star-forming regions described in a forthcoming study (Trayford et al., in prep.). However, here we use the ‘intrinsic’ (i.e. rest-frame and dust-free) colours of galaxies to examine the changes arising purely from the evolution of their stellar content. To simplify the interpretation we always quote rest-frame colours: there is therefore no ‘k’-correction needed to compare galaxies in the same band at different redshifts. We concentrate here on $u^*r^*$ colours (with the * referring to intrinsic colours) rather than $g^*r^*$, because the $u$ band is more sensitive to recent star formation, leading to more clearly separated blue/red colour sequences. Indeed, the $u^*-r^*$ index traverses the 4000Å break, often used as a proxy for star formation activity (e.g. Kauffmann et al. 2003). The photometry is presented here without dust effects, comparison is possible with various observational data where dust corrections have been estimated (e.g. Schawinski et al. 2014).

3 COLOUR EVOLUTION OF THE ENSEMBLE GALAXY POPULATION

Figure 1a shows that a scatter plot of EAGLE galaxies in a colour-stellar mass diagram, $(u^*-r^*)$ vs $M_\star$, exhibits strong bimodality in colour at redshift $z \approx 0$. The well defined red sequence resides at $u^*-r^* \gtrsim 2.2$ with colours becoming redder with increasing $M_\star$. The blue cloud is at $u^*-r^* \approx 1.3$, with a slope similar to that of the red sequence. These two sequences are indicated by red and blue lines to guide the eye, respectively, obtained by a spline fit to the maxima in the probability distribution of $u^*-r^*$ in bins of $M_\star$. We keep the location of these lines fixed in Fig. 1b–d to facilitate comparison at higher $z$. We clearly see that:

(i) The red sequence becomes bluer and less populated with increasing $z$. It is in place at $z \approx 1$ but has mostly disappeared by $z \approx 2$. A gap in the red sequence is noticeable at $z \approx 1$ for $M_\star \sim 10^{10.7} \, M_\odot$.

(ii) The blue sequence becomes bluer with increasing $z$, and exhibits decreasing scatter.

The main features of the galaxy population that drive these trends are illustrated in the bottom panels of the figure. Figs. 1e & g show that for a narrow stellar mass range around $M_\star = 10^{10.25} \, M_\odot$ that $u^*-r^*$ is strongly anti-correlated with the specific star formation rate, sSFR $\equiv M_\star/\dot{M}_\star$, provided the galaxy is star-forming ($\log_{10}(\text{sSFR/Gyr}) \gtrsim -2$). This is not surprising since the light in the $u^*$ filter is dominated by emission from massive and hence young stars, while $r^*$ is dominated by the older population. Galaxies in this plot follow a very tight relation at a given $z$, sliding along a narrow locus in colour that becomes bluer at higher $z$. At $z = 2$, the galaxies follow almost the same relation in $u^*-r^*$ versus sSFR as at $z = 0$, with just a small but noticeable offset towards redder colours ($\approx 0.1$ mag) for $\log_{10}(\text{sSFR/Gyr}) \lesssim -1.25$. This is a result of the redder population being younger on average and hence brighter for a $z = 2$ star-forming galaxy, compared to a star-forming galaxy at $z = 0$. For lower star formation rates ($\log_{10}(\text{sSFR/Gyr}) \lesssim -1.25$) the old population has more influence on $u^*-r^*$, and the younger average stellar age of $z = 2$ galaxies causes an offset to blue colours.

With $u^*-r^*$ colour so strongly correlated with sSFR, the colour vs $M_\star$ diagram of Fig. 1 is one view of the fundamental plane of star-forming galaxies, discussed recently by Lagos et al. (2015a). These authors showed that EAGLE galaxies from different redshifts fall onto a single 2D surface when plotted in the 3D space of $M_\star - M_\star$ and gas fraction (or metallicity), which they attributed to self-regulation of star formation. Lagos et al. (2015a) also showed that observed galaxies follow very similar trends. The increasingly bluer colours of the blue cloud towards higher $z$ is a consequence of the increased star formation activity at fixed $M_\star$. 

Downloaded from http://mnras.oxfordjournals.org/ at Liverpool John Moores University on June 23, 2016
The scatter in colour at fixed \( M_\star \) on the red sequence is mostly due to metallicity, \( Z \), as is clear from examination of the \( u^*-r^* \) distribution of galaxies at a given \( M_\star \), with low sSFR < \( 10^{-3} \) Gyr\(^{-1} \), plotted as green points in Fig. 1f. The colour of star-forming galaxies with sSFR > \( 10^{-3} \) Gyr\(^{-1} \) (black points) also depends on \( Z \), but from comparison of these panels it is clear that this effect is much smaller than the dependence of colour on sSFR itself - it induces the small scatter in \( u^*-r^* \) in panel e, on top of the main trend with sSFR.

As discussed by many others, the slope of the red sequence demonstrates the dependence of colour on \( Z \) for galaxies: more massive galaxies are more metal rich and hence redder (see Trayford et al. 2015). Because the mass- and light-weighted metallicities are not equivalent, the internal metallicity distribution for stellar populations in a galaxy may also affect the normalisation of the red sequence. To illustrate this, we calculated the median metallicity, \( Z_{\text{med}}(M_\star) \), in bins of stellar mass. We then calculated \( u^*-r^* \) colours for a 10 Gyr old population with that dependence of \( Z \) on \( M_\star \), and plot the resulting \( u^*-r^* \) colour as a function of \( M_\star \) in Fig. 1a as red squares. Although this sequence has the same slope as the red sequence in EAGLE, it is systematically redder by \( \approx 0.25 \) magnitudes. This is not an age effect, but a consequence of stellar populations exhibiting a spread in metallicity within an EAGLE galaxy. In fact, the metallicity distribution function of stars in an EAGLE galaxy is fairly well described by an exponential distribution.

The consistent red sequence slope between the two grey lines in panel d) at redshift \( z = 0.1 \) and \( z = 2 \). The olive line indicates the median \( u^*-r^* \) as a function of sSFR at \( z = 0.1 \) for comparison at \( z = 2 \). Panel f: \( u^*-r^* \) versus median stellar metallicity for the galaxies of panel e; galaxies with sSFR< \( 10^{-3.5} \) Gyr\(^{-1} \) appearing in the green box in panel e, are plotted as green dots.

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The consistent red sequence slope between the toy model and EAGLE suggests that any changes in the internal stellar \( Z \) distribution of EAGLE galaxies with mass are not strong enough to bias the median colours of red galaxies. We note that the slope of \( u^*-r^* \) relation as a function of \( M_\star \) in the blue cloud is set by the sSFR-\( M_\star \) relation and
not by metallicity effects. Therefore, the similarity between the slopes of the blue and red lines is coincidental.

3.1 Satellite colours
The extent to which satellite galaxies are preferentially red relative to the general population is illustrated in Fig. 2. We divide colour vs mass diagrams at different redshifts into equal bins of $u^* - r^*$ and $\log_{10}(M_*/M_\odot)$. The satellite fraction in each bin is computed and normalised by the total satellite fraction for all galaxies in the same stellar mass range. Bins containing $>10$ galaxies are shaded by the $\log_{10}$ normalised satellite fraction, such that positive values indicate a higher than average satellite fraction for that mass, while negative values indicate a lower than average value. The satellite fraction as a function of stellar mass in eagle is plotted for each redshift in the bottom panels.

Galaxies with $M_* \leq 10^{10}M_\odot$ that are red are predominately satellites, seen most strikingly at $z \approx 2$. At lower $z$ there is still a trend for low-mass red galaxies to be satellites, but the trend is less pronounced because some galaxies classified as centrals are also red. To some extent this may be a consequence of galaxies being quenched by ram-pressure stripping in the outskirts of more massive halos, before they are classified as being a satellite (e.g. Bahé et al. 2013). Indeed, they may not be part of the FOF halo (yet). Another possibility is that some of these galaxies were stripped as satellites when they fell inside a massive halo but have travelled out again, the so-called backsplash population (Balogh et al. 2000).

At redshift $z \approx 0$, the fraction of satellites is $\approx 50$ per cent at $M_* \sim 10^{10}M_\odot$, decreasing slowly to 30 per cent by $M_* \sim 10^{10.5}M_\odot$, and then dropping rapidly towards higher $M_*$. Satellite fractions decrease slowly at all $M_*$ with increasing $z$ to $z \approx 1$, and then drop much faster to below 30 per cent at all masses by $z = 2$. This rapid drop in the satellite fraction with increasing $z$ is the reason that the red sequence disappears at low $M_* \leq 10^{10}M_\odot$ for $z \gtrsim 2$.

3.2 AGN host colours
The effect of feedback from accreting black holes on galaxy colours is illustrated in Fig. 3. We only plot central galaxies to disentangle satellite quenching from effects induced by AGN. The figure is analogous to Fig. 2, with median black hole mass replacing satellite fraction. The median black hole mass ($M_\bullet$) as a function of $M_*$ is plotted for each redshift as the bottom panels in Fig. 3.

At redshift $z = 2$, there is a very strong trend for red galaxies with $M_* \approx 10^{10}M_\odot$ to exhibit unusually high black
Figure 3. The impact of black hole mass on the evolution of the $u^*-r^*$ vs $M_*$ colour-stellar mass relation for central galaxies. Top panels: Each square corresponds to a bin in $u^*-r^*$ and $M_*$, and is coloured according to the median black hole mass, $M_\bullet$, in that bin, such that larger values of $M_\bullet$ correspond to redder colours (see the colour bar). The median black hole mass in each square is normalised to the median black hole at that stellar mass (bottom panel), removing trends of $M_\bullet$ with $M_*$ and redshift. At $z = 2$ there is a trend for redder galaxies to have more massive black holes. This trend is particularly striking for galaxies with $M_* \sim 10^{10} M_\odot$ and becomes less pronounced at higher masses. There is no obvious correlation at lower stellar mass. These trends persist to lower $z$ but become weaker.

Lower panels: The median black hole mass, $M_\bullet$, as a function of stellar mass is plotted in solid black. Dashed black lines represent the 16th and 84th percentiles. $M_\bullet$ is nearly independent of $M_*$ below $M_* \sim 10^{10} M_\odot$, and increases with $M_*$ above this characteristic mass. This trend is almost independent of redshift.

3.3 Colour transformation mechanisms

Combining the results of Figs. 2 and 3 enables us to understand the origin of the evolution in the $u^*-r^*$ vs $M_*$ diagram of Fig. 1: galaxies with $M_* < 10^{10} M_\odot$ tend to become red when they become satellites, whereas galaxies above this characteristic mass are quenched by their AGN. This reasoning also explains why the red sequence starts to build-up from both the low-mass and the high-mass ends, leaving initially a noticeable scarcity of red galaxies at $M_* \approx 10^{9.7} M_\odot$ at $z \approx 1$. Such galaxies are too low mass to host a vigorously accreting black hole, yet too massive to be satellites in the typically lower-mass groups at that higher $z$. It is not until redshifts $z < 1$ that the more massive halos that host $M_* \approx 10^{10} M_\odot$ satellites appear.

The extent to which EAGLE predicts the characteristic stellar mass above which AGN quenching occurs, and the evolution of the under-abundance of intermediate-mass red galaxies, not only depends on the details of the sub-grid physics but also on the volume that is simulated. This is because massive clusters are under-represented or simply absent due to missing large-scale power in the density field, and poor sampling of rare objects in the relatively small, periodic EAGLE volume of $100^3$ cMpc$^3$. However we believe that the relevant physics described here is robust, and cor-
4 COLOUR EVOLUTION OF INDIVIDUAL GALAXIES

4.1 The flow of galaxies in the colour- \( M_\star \) plane

Fig. 4 illustrates how galaxies move through the \((u^* - r^*) - M_\star\) plane. Selecting galaxies in equal bins of \(u^* - r^*\) and \(\log_{10}(M_\star)\) at one redshift, we measure the median difference in \(u^* - r^*\) and \(\log_{10}(M_\star)\) for their descendant galaxies at a second redshift. We plot these differences for galaxies over an equal time period at high \((z \approx 1)\) and low \((z \approx 0)\) redshift. This is achieved by using two consecutive snapshots \((z = 1.3\) and \(z = 1)\) for the right panel, corresponding to a time interval of \(\approx 0.9\) Gyr, and interpolating the galaxy vectors between the lowest redshift snapshots \((z = 0.1\) and \(z = 0)) to match the same time period. Descendants that grow in \(M_\star\) by a factor > 4 through merging into a more massive host galaxy are eliminated from the measurement, to prevent them contributing extreme vectors to their bin.

For both redshift ranges, it is clear that the colours of galaxies generally become redder, with vectors pointing in the positive \(u^* - r^*\) direction. Exceptions can be seen on the red sequence, in which some red galaxies become star-forming following a gas-rich merger - we discuss the fraction of such 'rejuvenated' galaxies below. Red sequence galaxies show little change in \(u^* - r^*\), but in general those with \(M_\star \lesssim 10^{10.75} M_\odot\) lose mass, with only the most massive red sequence galaxies showing mass growth. Considering that we do not count mergers into hosts of factor > 4 higher \(M_\star\), this suggests that red sequence galaxies are being stripped prior to a dry merger with a massive central. We also see evidence of 'mass quenching', with the vectors for blue-cloud selected galaxies becoming steeper with increasing stellar mass. For \(M_\star > 10^{10} M_\odot\) this is attributable to the presence of AGN (Fig. 3).

Another notable behaviour seen in Fig. 4 is that the bluest galaxies \((u^* - r^* \approx 0.5,\) right panel\) tend to change their \(u^* - r^*\) colour more than the average blue galaxy, to the extent that they end-up on the red side of the blue sequence at the later redshift. The strong reddening and mass increase of these galaxies suggests that they are starbursts triggered just prior to a merger. Such a scenario is consistent with the fact that a higher proportion of satellites are found at these colours than at the centre of the blue peak (see Fig. 2). Note that this effect is only observed for the high redshift panel, partially due to the bluest bins failing to meet the minimum galaxy count criterion of 10. These extreme starbursts are clearly rarer at low redshift.

4.2 Evolution of colour populations in EAGLE

To track the evolution of galaxies selected to be red, blue or green, we must first define these populations. To do this we apply cuts that evolve with redshift \(z\) for red and blue galaxies.

\[
(u^* - r^*)_{\text{red}} > 0.2 \log_{10}(M_\star/M_\odot) - 0.25 z^{0.6} + 0.24
\]
\[
(u^* - r^*)_{\text{blue}} < 0.2 \log_{10}(M_\star/M_\odot) - 0.25 z^{0.6} - 0.3(1)
\]

The green galaxies are taken to be those that are not included in either set. These cuts are defined in an \textit{ad-hoc} way to divide the galaxies into three populations at each redshift. This is a similar procedure to that used by many observational studies, with authors adopting differing functional forms and normalisations (see e.g. the discussion in Taylor et al. 2015). The exact form of the colour cuts is unimportant for our qualitative analysis, but is considered when we discuss our quantitative results.

The evolution of the \(u^* - r^*\) colours of galaxies, selected by colour either at high redshift \((z > 0.5)\) or low redshift \((z < 0.1)\), is illustrated in Fig. 5. We bin galaxies in the three colour bins described above and plot the colour distribution of the descendants and main progenitors (odd and even rows, respectively) of galaxies selected to be red, green or blue (top two, middle two, and bottom two rows, respectively). We use the galaxy merger trees to identify descendants and main progenitors. For each panel, the \(u^* - r^*\) colour distribution of all galaxies at the indicated redshift, with \(M_\star > 10^{10} M_\odot\), is plotted in grey.

From the top two rows it becomes clear that most galaxies that are red at \(z = 0.5\) stay red down to \(z = 0\), whereas a substantial fraction of galaxies that are red at \(z = 0\) were green at \(z = 0.1\) or even blue at \(z = 0.5\). Galaxies that are green at \(z = 0\) predominantly become red at \(z = 0\), but a fraction of green galaxies becomes blue (third row). Galaxies that are green at \(z = 0\) had a range of colours at \(z = 0.5\), although they were bluer than average (fourth row). The similar and dominant blue (red) fractions in the two green progenitor (descendant) panels suggests that the typical time to transition through the green valley is shorter than the redshift intervals used here. Finally, galaxies that are blue at \(z = 0.5\) have a large range of colours at \(z = 0\) with a distribution that is similar to that of the population as a whole (fifth row), whereas galaxies that are blue at \(z = 0\) were mostly blue at \(z = 0.5\) as well (bottom row).

The rate at which galaxies with stellar mass \(M_\star > \)
10^{10} M_\odot at z = 0.5 change u'-r' colour over the redshift range z = 0.5 to z = 0 (elapsed time \Delta t \approx 5 \text Gyr) is quantified in Fig. 6. We identify the z = 0 descendant for all galaxies with \( M_* > 10^{10} \text M_\odot \) at \( z = 0.5 \), compute the change in colour, \( \Delta(u'-r') \), and plot a histogram of rates, \( \Delta(u'-r')/\Delta t \). We also identify if a galaxy is a satellite at \( z = 0.5 \), or if the mass of its central black hole increases by a factor \( \geq 1.5 \). This threshold is chosen to represent an above average black hole growth, while still providing a significant sample of galaxies.

The rate of change of the median colour of galaxies is small, \( \Delta(u'-r')/\Delta t \approx 0.08 \text mag \text Gyr}^{-1} \) to the red, but is larger for galaxies whose black hole grows more than average (\( \Delta(u'-r')/\Delta t \approx 0.09 \text mag \text Gyr}^{-1} \)) or those that are satellites (\( \Delta(u'-r')/\Delta t \approx 0.12 \text mag \text Gyr}^{-1} \)). Galaxies that are red at \( z = 0.5 \) typically change little in colour to \( z = 0 \), (\( \Delta(u'-r')/\Delta t \approx 0.03 \text mag \text Gyr}^{-1} \)), except for the occasional outlier that becomes blue. The rate of change of the median colour is larger for galaxies that are green or blue at \( z = 0.5 \), with individual galaxies changing colour more rapidly, both to the red and to the blue. Galaxies that are satellites can undergo rapid changes to the red, \( \Delta(u'-r')/\Delta t \approx 0.2 \text mag \text Gyr}^{-1} \), whether blue or green at \( z = 0.5 \). Note that this rate is averaged over a considerable period (\( \approx 5 \text Gyr \)), and instantaneous rates of colour change for galaxies can be much higher, as explored below.

### 4.3 Colour-mass tracks of individual galaxies

We have examined a large number of tracks of individual galaxies in (\( u'-r', M_* \)) space and have identified three generic tracks of central galaxies that we illustrate in Fig. 7. In Fig. 8 we also show the track of a central galaxy that is very massive at \( z = 0 \) (\( M_* \sim 10^{11} \text M_\odot \)), to illustrate individual tracks of satellites that merge with it. More details of the four galaxies tracked in these panels are given in Table 1. The time-scale over which galaxies transition to the red sequence is compared to that of a passively evolving population, by plotting (single) star-burst tracks initiated at different times (grey curves in middle panels).

The blue track in Fig. 7 is for a galaxy that remains in the blue cloud down to \( z = 0 \), forming stars in a blue disc that grows in time. As its sSFR decreases with time, and the contribution of an older population of stars becomes more important, it slowly reddens, with \( \Delta(u'-r') \approx 0.2 \) from \( z = 1 \) to \( z = 0 \) (elapsed time \( \approx 8 \text Gyr \)). We see from Table 1

![Figure 4](image_url)  
**Figure 4.** The flow of galaxies in (\( u'-r', M_* \)) space between two redshifts, \( z_1 \) to \( z_2 \). Right panel shows the galaxy flow between \( z_1 = 1.3 \) and \( z_2 = 1 \) snapshots, a period of \( \approx 0.9 \text Gyr \). The left panel shows the galaxy flow interpolated between the \( z = 0.1 \) and \( z = 0 \) snapshots to yield the same time period, with \( z_1 = 0.07 \) and \( z_2 = 0 \). Black circles represent the mean location of galaxies at \( z_1 \), selected in a bin of \( u'-r', M_* \); the size of the circle is proportional to the logarithm of the total stellar mass in galaxies in that bin. Black vectors represent the mean motion of the galaxies in that bin between \( z_1 \) and \( z_2 \). Orange vectors (purple vectors) are for those galaxies that at redshift \( z_2 \) belong to halos with virial mass \( M_{200, \text{crit}} < 10^{13} \text M_\odot \) (\( M_{200, \text{crit}} > 10^{13} \text M_\odot \)). Centre coordinates and vectors sampling fewer than 10 galaxies are not plotted. The overall distribution at the later redshift is plotted as grey contours for comparison. We do not take into account galaxies merging into hosts that are more than four times their mass, illustrating such mergers in more detail in Fig. 8.

### Table 1. Properties of the galaxies plotted as main tracks in Figs. 7-8.

| Sym./Fig. | GalaxyID | \( M_*/\text{M}_\odot \) | Satellite |
|-----------|----------|----------------|-----------|
| Circle/7  | 18169630 | \( 7.09 \times 10^6 \) | n         |
| Square/7  | 15829793 | \( 1.03 \times 10^6 \) | n         |
| Triangle/7| 1496270  | \( 6.00 \times 10^7 \) | n         |
| Circle/8  | 15197399 | \( 1.84 \times 10^8 \) | y         |
Figure 5. Grey histograms: $u^*-r^*$ colour evolution of all galaxies from redshift $z = 0.5$ to $z = 0$, with redshift decreasing from left to right in each row; each panel is labelled with the corresponding redshift. Only galaxies with $M_* > 10^{10} M_\odot$ are included. Colour selections are made using Eq. 1 and are as follows: top row: we select red galaxies at $z = 0.5$ (red histogram) and plot the colour distribution of their descendants at low $z$ as a red histogram. Second row from top: we select red galaxies at $z = 0$ (red histogram), and plot the colour distribution of their main progenitors as a red histogram at higher $z$. Rows 3 and 4 from the top: as above, but for green galaxies. Bottom two rows: as above, but for blue galaxies. The background colour of the panel in which galaxies were selected is coloured grey for ease of reference.

and Fig. 3 that the black hole mass is $\approx 0.6$ dex lower than the median value for its stellar mass (at approximately the 3rd percentile of galaxies for that $M_*$), suggesting low levels of black hole feedback in the galaxy’s history.

The red track in Fig. 7 corresponds to a galaxy that becomes red more rapidly, reddening by $\Delta(u^*-r^*) \approx 1$ in $\approx 2$ Gyr, joining the red sequence at $z = 2$. From then on, its colour or stellar mass hardly changes; it has been a very compact elliptical galaxy since at least $z = 1$, maintaining a stellar half-mass radius of $< 2.3$ pkpc. The similar rate of colour transition for this galaxy to that of an instantaneous starburst suggests rapid quenching of star formation. Considering the black hole mass in Table 1 and the bottom-left panel Fig. 3 shows that the galaxy has a central black
Figure 6. Top left panel: Distribution of the average rate of colour change over the redshift interval $z = 0.5 \rightarrow 0$, $\Delta (u^*-r^*)/\Delta t$, for galaxies with mass $M_\star > 10^{10} M_\odot$; vertical arrows (top and bottom) denote the median rate. Other panels: As the top left panel, but for galaxies selected at $z = 0.5$ to be blue, red, and green (clockwise from top right). In each panel, black lines refer to all selected galaxies, purple lines to the fraction that at $z = 0.5$ are satellites, orange lines to the fraction whose central black hole has grown by at least a factor of 1.5 between $z = 0.5$ and $z = 0$. The rate of change of the median colour is typically small, but individual galaxies can change colour more dramatically over this period, $-0.2 < \Delta (u^*-r^*) < 0.25$ up to $|\Delta (u^*-r^*)/\Delta t| = 0.25$ Gyr$^{-1}$, particularly for green and blue galaxies.

The thick green track shows a galaxy that reddens at a similar rate ($\Delta (u^*-r^*) \approx 1.5$ in $\approx 2$ Gyr) at $z \approx 1$, after undergoing a near-equal mass merger (thin green line indicates the track of the other galaxy) leading to significant AGN growth. It then changes morphology from being a disturbed disc to a compact elliptical at $z = 0.5$. At redshift $z = 0.25$, it starts forming stars again, turning blue ($\Delta (u^*-r^*) \approx -1$) over $\approx 2$ Gyr. By $z = 0$, it has grown significantly in size, with a prominent bulge and an extended distribution of stars around it. The value of $M_\star$ in Table 1 puts the black hole mass at approximately the 75th percentile of galaxies for that stellar mass.

Note that the black hole masses for galaxies selected in Fig. 7 (and listed in Table 1) are all in the upper or lower quartiles for their stellar mass. While galaxies with average black hole masses have diverse histories, more extreme $M_\star$ values can be indicative of the certain types of track represented here. The quiescently star forming and quenched galaxies are both associated with particularly low and high levels of black hole growth respectively. The rejuvenated galaxy has a high black hole mass, though less extreme. The black hole growth is associated with a quenching event at $z \approx 1$, but the galaxy re-accretes gas and is able to recommence prolonged ($> 2$ Gyr) star formation in spite of its...
high black hole mass. We saw in Fig. 3 that for $z = 0$ and at these masses the overall correlation between galaxy colour and the residuals of the $M_\star - M_{\ast}$ relation is rather weak.

The massive $z = 0$ galaxy in Fig. 8 is a blue star-forming disc until just below $z = 1$, after which it becomes red and evolves into an elongated elliptical. Thin lines show the tracks of five galaxies that merge with it, with the line colour changing from purple to orange while these galaxies become satellites. Examination of these tracks reveals that while some galaxies become red when they are still centrals, most galaxies quench after being identified as satellites. This suggests that satellite identification is a good predictor of...
colour change in EAGLE, particularly for galaxies falling into a more massive halo. The satellite tracks exhibit a characteristic shape of rapid quenching followed by stellar mass loss, as the galaxy is stripped and eventually merges. We also see that the central galaxy exhibits quite a stochastic colour evolution compared to those in Fig. 7. This is perhaps due to the higher frequency of satellite interactions and mergers for the higher mass halo represented in Fig. 8, with only the rejuvenated track of Fig. 7 showing (two) mergers with satellites of $M_\ast > 10^9 M_\odot$.

The heterogeneous colour evolution of galaxies illustrated in Fig. 7 and 8 are consistent with observations that the green valley population is diverse (e.g. Cortese & Hughes 2009; Schawinski et al. 2014). In particular, Cortese & Hughes (2009) also find examples of galaxies moving off the red sequence through re-accretion of gas reminiscent of the green track in Fig. 7.

To quantify the fraction of galaxies that undergo a rapid colour transformation, we trace the main progenitors of all galaxies with mass $M_\ast > 10^{10} M_\odot$ at $z = 0$ back in time (a sample of ≈ 3000 galaxies). We find that the fraction of galaxies with at least one $\Delta(u^* - r^*) > 0.8$ change over any 2 Gyr period in their history is ≈ 40 per cent. We will refer to such galaxies as rapidly reddening galaxies; they tend to follow a track similar to the red track in Fig. 7. Of the galaxies that redden quickly, ≈ 1.6 per cent undergo a colour change to the blue of $\Delta(u^* - r^*) < -0.8$ over a 2 Gyr period. We will refer to this small fraction of galaxies as ‘rejuvenated’; they
follow a track similar to the green track in Fig. 7. Galaxies that do not ever undergo such a rapid reddening event make-up 60 per cent of the sample. We will refer to these as ‘quiescently star-forming’ galaxies; they follow a track similar to the blue track in Fig. 7.

Fig. 9 shows the distribution of $z = 0$ colours for $10^{10} < M_* / M_\odot < 10^{10.5}$ galaxies classified as having undergone a rapid transformation to redder colour (orange histogram), and those that never underwent such a rapid reddening (quiescently star-forming galaxies, purple histogram). A small fraction of galaxies become red without ever experiencing a rapid reddening event: this is the tail of the purple histogram towards red $u^*-r^*$ colour. Similarly, there is a tail to blue $u^*-r^*$ colour in the orange histogram, representing galaxies that underwent a rapid colour transition to the red, followed by more recent star formation turning them blue once more. The fraction of these galaxies is much higher than the 1.6 per cent of galaxies we classified as ‘rejuvenated’: the majority of galaxies that are blue now but were red in the past ($\approx 10\%$ of total), became blue more gradually than the $\Delta(u^*-r^*) = -0.8$ over 2 Gyr that we used to define ‘rejuvenated galaxies’.

Galaxies must become blue rapidly or traverse the entire green valley to meet these criteria. However, it is also interesting to note the probability that a galaxy identified in the green valley is on a bluer trajectory in the colour-$M_*$ plane. We identify this for green valley galaxies at $z < 2$, enforcing that a galaxy must undergo a monotonic colour change of $\Delta(u^*-r^*) < -0.05$ to be deemed significant (above the level of photometric error, e.g. Padmanabhan et al. 2008). We find that the EAGLE green valley galaxies have a 17% chance of being on a significantly blue trajectory. Conversely, we find the probability of a green galaxy being on a significantly red trajectory to be 75%, taking $\Delta(u^*-r^*) > 0.05$ as the criteria for being significant.

The tracks of individual galaxies enable us to characterise the colour transition time-scale of a galaxy by the time interval, $\Delta t_{\text{green}}$, it spent in the green valley on its way from the blue cloud to the red sequence. We calculate $\Delta t_{\text{green}}$ as follows: using Eq. (1) we select red galaxies at $z = 0$ and trace their main progenitors back in time to identify the last time ($t_1$) they became red and the last time ($t_2$) they were blue. Histograms of colour transition times, $\Delta t_{\text{green}} = t_1 - t_2$, for galaxies that at $z = 0$ have $10^{10} < M_*/M_\odot < 10^{10.5}$ and $10^{10.5} < M_*/M_\odot < 10^{11}$ are plotted in Fig. 10 (top and bottom panels, respectively).

The mode of the colour transition time distribution is $\approx 1.5$ Gyr, with a median of $\approx 2$ Gyr, mostly independent of whether quenching is likely due to becoming a satellite or AGN activity (purple and orange histograms, respectively). This is the time-scale for a passively evolving blue population of stars to turn red, as can be seen from Fig. 7. Strikingly, there is a very long tail to high values in the distribution of $\Delta t_{\text{green}}$ as inefficient quenching allows a small fraction of galaxies to spend a long time in the green valley before eventually turning red, whether due to becoming a satellite or hosting an AGN. Though quenched galaxies are more prevalent in high-mass halos, the colour transition time-scales show little dependence on halo mass. Despite this, the longest time-scales we measure are for halo masses $< 10^{13} M_\odot$.

The time-scales for colour transition and the quenching of star formation are clearly linked, however colour transition times are longer due to the passive evolution of stellar populations. This is illustrated by the grey curves in Fig. 7 and 8, showing that the colour transition time for an SSP is $\approx 2$ Gyr. The quenching time-scales of observed satellites presented by e.g. Muzzin et al. (2014) and Wetzel et al. (2013) are significantly shorter, typically $\lesssim 0.5$ Gyr and $\lesssim 0.8$ Gyr respectively. Although the $u^*-r^*$ colour index alone is not sensitive enough to resolve these quenching times, the typical colour transition times of EAGLE galaxies are consistent with such rapid quenching.

However, there are cluster studies in the $0 < z < 0.5$ redshift range that infer longer timescales ($\gtrsim 1$ Gyr, e.g. von der Linden et al. 2010; Vulcani et al. 2010; Haines et al. 2013). Fig. 10 does show a tail to longer quenching times for both AGN and satellites, though they are not typical at low redshift. Ultimately, the different observational tracers and modelling used to infer quenching times are still subject to significant systematics that may explain the different measurements (McGee et al. 2014). In particular, colour transition and quenching timescales are not equivalent, and galaxies may move between the red and blue populations when observed in optical and UV colour (Cortese 2012, e.g.). An analysis of the physical quenching timescale and its evolution is left to a future study.

5 CONCLUSIONS

We have investigated the evolution and origin of the colours of galaxies in the EAGLE cosmological hydrodynamical simulation (Schaye et al. 2015; Crain et al. 2015). We apply the single population synthesis models from Bruzual & Charlot (2003) to model galaxy colours in the absence of dust, as described by Trayford et al. (2015). We also use galaxy merger trees to trace descendants as well as main progenitors through time.

The $u^*-r^*$ vs $M_*$ diagram is bimodal at redshift $z = 0$, with a clearly defined red sequence of quenched galaxies, and a blue cloud of star-forming galaxies. The scatter and slope of the red sequence are both determined mainly by stellar metallicity, while the normalisation additionally depends on stellar age (Fig. 1). The scatter in the blue cloud, in contrast, is mostly due to scatter in the specific star formation rate at fixed stellar mass (sSFR $= M_\star / M_\odot$). The slope of blue cloud colours versus $M_\star$ is similar to that of the red sequence, but as their origins are different, this is coincidental. At higher $z$, both colour sequences become bluer, and the red sequence becomes less populated until it has mostly disappeared by $z = 2$.

From studying the evolution of EAGLE galaxies in $u^*-r^*$ and $M_\star$, we note that in general:

- Galaxies in EAGLE turn red either because they become satellites (mainly at lower masses, see Fig. 2) or because of feedback from their central supermassive black hole (mainly

2 Galaxies may also undergo slower colour transitions due to secular evolution, and the number density of galaxies does not remain constant because of mergers. These quoted fractions therefore inevitably depend on how galaxies are selected, i.e. the colour choice we made in Eq. 1.
for more massive galaxies, see Fig. 3). As a consequence, the red sequence builds-up from both the low-mass and high-mass sides simultaneously, with the low-mass red galaxies being satellites of the massive red centrals that are quenched by their AGN. This results in a dearth of red galaxies at intermediate mass, \( M_* \sim 10^{10} \, M_\odot \), at \( z \approx 1 \) - such galaxies are too low mass to host a massive black hole, but too massive for a large fraction of them to be satellites in the EAGLE volume. While we believe the existence of such a deficit is unlikely to change with increased simulation volume, it should be noted that the limited volume and lack of large scale power in the EAGLE 100^3 Mpc\(^3\) simulation may affect the depth of the deficit.

- The colour evolution in the blue cloud is driven by the decrease in the sSFR rates of star-forming galaxies with cosmic time.

- The characteristic time scale for galaxies to cross the green valley, from the blue cloud to the red sequence (Fig. 10), is \( \Delta \text{green} \approx 2 \, \text{Gyr} \), mostly independent of galaxy mass and cause of the quenching. It is determined by the rate of evolution of a passive population of blue stars to the red. This timescale is consistent with rapid or instantaneous quenching of star formation, as inferred from observations of satellite galaxies by Muzzin et al. (2014). The distribution of \( \Delta \text{green} \) has an extended tail to \( \approx 10 \, \text{Gyr} \): a small fraction of galaxies spend only a short time, \( \Delta \text{green} \lesssim 2 \, \text{Gyr} \), in the green valley - it is not easy being green.

We identified three characteristic tracks that galaxies follow in the \( u^*-r^* \) vs \( M_* \) diagram (Figs. 7 and 8). Quiescently star-forming galaxies remain in the blue cloud at all times, without sudden reddening episodes of \( \Delta (u^*-r^*) > 0.8 \) in any 2 Gyr interval. Nearly 60 per cent of galaxies with stellar mass at \( z = 0 \) greater than \( M_* = 10^{10} \, M_\odot \) fall into this category (see Fig. 9). The remaining 40 per cent of galaxies do undergo such sudden episodes of star formation suppression. The majority of these rapidly reddened galaxies move onto the red sequence permanently as per the evolutionary picture of Faber et al. (2007), however we find that 1.6 per cent undergo an episode in which star formation causes the galaxy to change colour to the blue again, having \( \Delta (u^*-r^*) < -0.8 \) over a 2 Gyr period (e.g. Fig. 7). The fraction of such rejuvenated galaxies is thus very small. Nevertheless, a much larger fraction of the galaxies that at \( z = 0 \) are blue were red in the past: the rate of colour transition of galaxies to the blue is generally significantly slower than the quenching timescale. We also find that the fraction of green valley galaxies on blue trajectories (where \( \Delta (u^*-r^*) < -0.05 \)) at a given instance from \( z < 2 \) is larger still at 17%, implying that only a subset transition completely from red to blue and remain there.

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