Properties of cell signaling pathways and gene expression systems operating far from steady-state

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Ligand-receptor systems, covalent modification cycles, and transcriptional networks are basic units of signaling systems and their steady-state properties are well understood. However, the behavior of such systems before steady-state is poorly characterized. Here, we analyzed the properties of input-output curves for each of these systems as they approach steady-state. In ligand-receptor systems, the EC\(_{50}\) (concentration of the ligand that occupies 50% of the receptors) is higher before the system reaches steady-state. Based on this behavior, we have previously defined PRESS (for pre-equilibrium sensing and signaling), a general "systems level" mechanism cells may use to overcome input saturation. Originally, we showed that, given a step stimulation, PRESS operates when the kinetics of ligand-receptor binding are slower than the downstream signaling steps. Now, we show that, provided the input increases slowly, it is not essential for the ligand binding reaction itself to be slow. In addition, we demonstrate that covalent modification cycles and gene expression systems may also operate in PRESS mode. Thus, nearly all biochemical processes may operate in PRESS mode, suggesting that this mechanism may be ubiquitous in cell signaling systems.

Cells detect input signaling molecules using receptors, proteins usually located on the cell surface embedded in the plasma membrane. Activated receptors then transmit the signal to the interior of the cell through a series of downstream processes that typically lead to changes in gene expression, resulting in an appropriate output response to the input. In a way, the system's overall input-output curve summarizes its biological characteristics and function\(^1\). Several of the input-output features are well appreciated\(^2\). For example, some input-output curves are graded, changing outputs moderately over a wide range of inputs; others are ultrasensitive, changing very rapidly from low to high output across a narrow range of inputs. The response of biological systems to changing environmental conditions is a dynamic process. The response time (the time to reach 63% of its value at steady state) of a given system depends on the input strength and duration, as well as on the structure and the kinetic parameters of the system under consideration\(^3\). These two notions, the shape of an input-output curve and the dynamics involved in reaching its steady-state, have not been properly integrated in the literature. The focus of the current work is on studying how an input-output curve evolves over time and how this evolution confers dynamic tunability to the signaling system.

In a ligand-receptor system, the EC\(_{50}\) (the concentration of the ligand that occupies 50% of the receptors) changes over time\(^4\). In the case of a single binding site, the EC\(_{50}\) is elevated at early times, and it drops to its steady-state value as binding approaches equilibrium. At that time, the EC\(_{50}\) is equal to the dissociation constant of the ligand-receptor reaction. Thus, in effect, the dose-response binding curve shifts over time from right to left\(^4\). All these ideas, first stated in our previous article, are illustrated in Fig. 1. Note that a shift in the binding curve over time implies that a ligand–receptor system is sensitive (i.e., a change in the input concentration elicits a change in the output) in different regions of ligand concentration at different times before steady-state\(^4\). Based on this property, we observed that when the ligand-receptor complex activates a downstream signaling component...
with a time-scale compatible with acting before the ligand:receptor binding process achieves equilibrium, this signaling component will use the information contained in the pre-steady-state binding curve, and thus produce an output with an EC$_{50}$ shifted to the high ligand concentration region. Therefore, such a system provides distinguishable responses to ligand concentrations that saturate the receptors at steady-state. We termed this “systems level” mechanism PRESS, for Pre-Equilibrium Sensing and Signaling, and we showed its potential role in yeast directional polarization in response to pheromone gradients. Others have also shown the importance of pre-steady state information. For example, pathway dynamics has demonstrated to be essential to develop a prognostically useful model based on patient-stratification. Also, understanding time-dependent input-output curves can help understand mechanisms to confer transient bistability.

We have coined the term shift for the property of systems that change their EC$_{50}$ over time and refer to systems that exhibit this property as shifters. A shift mathematically means that the EC$_{50}$ has to be a time dependent function. In a ligand-receptor system, this shift is possible because the time to reach binding equilibrium depends inversely on the ligand concentration: the higher the ligand concentration, the faster the equilibrium is reached, see Fig. 1B,D. Thus, at early times after addition of the ligand, the binding process will be further from equilibrium at the lowest ligand concentrations, see Fig. 1C. This will result in a higher EC$_{50}$ value than at equilibrium and a steeper binding curve (i.e., less graded, Fig. 1C,E). Curve steepness is often related to the term ultrasensitivity as defined by Golbeter and Koshland, denoting a relationship between a change in input and a change in output. There are local and global definitions of sensitivity. Typically, an overall, global value for the ultrasensitivity of an input-output curve is taken as the EC$_{90}$/EC$_{10}$ ratio, where EC$_{90}$ and EC$_{10}$ are the inputs required to elicit 90% or 10% of the maximum output, respectively. For a Michaelian response, EC$_{90}$/EC$_{10}$ equals 8.1. Thus, the ratio log(81)/log(EC$_{90}$/EC$_{10}$), known as the Hill coefficient ($n_H$), is a global quantification of the curve steepness relative to that of a Michaelian response. This definition does not require a particular shape for the input-output curve.

Here, we asked if the concept of PRESS may be applied to processes that do not reach thermodynamic equilibrium, such as ligand-receptor binding, but do reach a steady-state, such as covalent modification (CM) cycles and gene expression systems. We found that both can indeed be shifters. Furthermore, we show that by introducing dynamics to the input it is possible to extend PRESS to fast signaling components, effectively extending this property to most signaling motifs and parameter sets.

**Figure 1.** Characterization of a prototypical shifter: a ligand-receptor reaction. (A) Ligand-Receptor scheme, ligand binds free receptor reversibly with rate constants $k_{on}$ and $k_{off}$. (B) Output (C, ligand-receptor complex) vs. time, for different values of input (color-coded using a heat-map scale). Circles over the curves mark output at time $\tau$ (time at which output is 63.2% of its steady-state value). (C) Input-output curves obtained at different times (color-coded using a heat-map scale). Circles mark the EC$_{50}$. (D) $\tau$ vs. input. (E) EC$_{50}$ (red) and $n_H$ (green) vs. time.
Results

Ligand-receptor systems revisited: a slowly increasing input controls the shift. Under physiological conditions, inputs (e.g., concentrations of growth factors or nutrients) are likely to change gradually rather than in a stepwise manner. Therefore, to determine how PRESS is altered when the input has dynamics, we considered the case of a ligand L that increases exponentially and interacts with a receptor R to form an activated complex LR, with binding and unbinding rates k1 and k2, respectively. We assumed that total amount of receptor is constant R0 = R + C, and that the concentration of free L is not depleted significantly by the binding reaction, so that L ~ LT (total amount of ligand). Thus, we described the system by the following equations:

\[ \frac{dy}{dt} = x(t_n)(1 - y) - y \]  
\[ x(t_n) = x_{\text{max}}(1 - \exp\left(-\frac{t_n}{\tau_{L,R}}\right)) \]

where \( x(t_n) = L(t_n)/K_D \) is the amount of ligand relative to the dissociation constant for the binding-unbinding reaction \( K_D = k_{2/2}/k_{1/1} \), \( y = C/R_{0} \) is the amount of ligand-receptor complex relative to the total amount of receptors, and \( t_n = t/t_{ref} \) is the time expressed in units of a reference time \( t_{ref} = 1/k_{off} \) (see Supplementary Information for a detailed derivation). We further assumed that the amount of ligand increases following an exponential function characterized by a maximum value \( L_{\text{max}} \) and a characteristic time \( \tau_{L,R} \), resulting in \( x_{\text{max}} = L_{\text{max}}/K_D \) and \( \tau_{L,R} = \tau_{L,R}/t_{ref} \) which connects the time-scale associated with the ligand accumulation \( (\tau_{L,R}) \) with that of the binding-unbinding process \( (t_{ref} = 1/k_{off}) \).

We solved equation (1a,1b) using three input dynamics, much faster (nearly a step increase), of the same order or much slower than the binding reaction timescale (\( \tau_{L,R} = 0.01, 1 \) and 100, respectively), for a range of inputs \( x_{\text{max}} \) (Fig. 2A–C). These simulations showed two important results. First, the response \( y \) reached steady-state faster when the input was larger (Fig. 2D–F), even though the time for the input \( x \) to reach \( x_{\text{max}} \) was modeled so that it was independent of its magnitude. Consequently, the \( y \) vs \( x_{\text{max}} \) plot shows a clear shift in the binding curve over time (Fig. 2G–I), with a lower bound of \( \text{EC}_{50} = 1 \). Second, the dynamics of the stimulus affected the speed of the shift: larger values of \( \tau_{L,R} \) resulted in slower shifts. That is, a slow rising input caused a slow shift (Fig. 2J), providing a longer time to a pathway downstream to operate in PRESS mode.

We also computed the Hill coefficient \( n_H \) for the three input dynamics. In contrast to what we observed for \( \text{EC}_{50} \), \( n_H \) dynamics was independent of the dynamics of the input (Fig. 2K). This result suggested that the \( n_H \) dynamics only depends on the ligand-receptor binding reaction. To explore this further, we derived analytical results in which we eliminated the dynamics of the binding reaction altogether. That is, the ligand-receptor complex was in quasi-steady state with the value of the input \( x \) (see Supplementary Information for details). In this situation, we obtained:

\[ y = \frac{x(t_n)}{1 + x(t_n)} \]

This equation is characterized by

\[ \text{EC}_{50}(t_n) = 1/(1 - \exp(-t_n/\tau_{L,R})) \]  \hspace{1cm} (3.a)
\[ n_H(t_n) = 1 \]  \hspace{1cm} (3.b)
\[ \tau(x_{\text{max}}) = \tau_{L,R} \left[ \frac{1 + 0.368x_{\text{max}}}{0.368(1 + x_{\text{max}})} \right] \]  \hspace{1cm} (3.c)

which show that the \( \text{EC}_{50} \) is a decreasing function of time (equation (3.a)), indicating that the shift is conserved, but that the \( n_H \) is time independent and equal to 1 (equation (3.b)). This confirms that \( n_H \) dynamics is conferred exclusively by the binding reaction itself. Equation (3.c) indicates, that for the convolved process, \( \tau \) is a decreasing function of the input.

In summary, in previous work we had noted that operation in PRESS mode required that the “sensing” ligand-binding reaction be slower than the downstream “signaling” component. In that study, we considered a step-like stimulation, here corresponding to \( \tau_{L,R} \ll 1 \). Here, our results indicate (Fig. 2) that the rising time of the input \( \tau_{L,R} \) controls the speed of the shift of the input-output curves. Thus, by introducing dynamics to the rise in the concentration of the input, it is not essential any more that the sensing component be particularly slow. It is enough that the stimulus rises slowly relative to the downstream component, even in cases where the ligand-receptor interaction is fast.

Covalent modification cycles. CM cycles consist of a substrate \( S \) that is activated into \( S^* \) by an enzyme through a post-translational modification such as phosphorylation and deactivated by another enzyme through the removal of that modification. Here we studied their behavior pre-steady state in response to the accumulation of the activating enzyme \( E_{a} \), with constant deactivating enzyme \( E_{d} \). We used a mass-action kinetics mechanistic description (i.e., considering SE complex formation, as detailed in the Supplementary Information) (Fig. 3A). For this system, we considered that the input is \( x = (k_{d}/k_{d})/(k_{d}/k_{d}) \), the ratio of the maximum velocities of the activating and deactivating enzymatic reactions, with catalytic rate constants \( k_{d} \) and \( k_{d} \); the output is \( y = S^*/S_{T} \), the fraction of active substrate; and the two relevant parameters are \( K_{a} = K_{a}/S_{T} \) and \( K_{d} = K_{d}/S_{T} \), the Michaelis-Menten
constants \((K_m = (d + k)/a)\) relative to the total amount of substrate. In what follows, the time is expressed in units of a reference time \(t_{\text{ref}} = S_T/(k_2 E_d)\), so that \(t_n = t/t_{\text{ref}}\). (For the simplified Michaelis-Menten description, see Supplementary Information).

We further assumed that the activating enzyme increases following an exponential function, for example due to its accumulation after its expression is induced. This function is characterized by a maximum value \(E_{a,\text{max}}\) and a characteristic time \(\tau_{E_{a,\text{max}}}\) resulting in:

\[
x(t_n) = x_{\text{max}}(1 - \exp(-t_n/\tau_{E_{a,\text{max}}}))
\]

where \(x_{\text{max}} = \frac{k_1 E_{a,\text{max}}}{k_2 E_d}\) and \(\tau_{E_{a,\text{max}}} = \tau_{E_{a,\text{max}}}/t_{\text{ref}}\).

Depending on the values of the Michaelis-Menten constants \(K_a\) and \(K_d\), two extreme regimes may be considered. In one regime, \(K_a\) and \(K_d\) are large \((K_m,a \gg S_f\) and \(K_m,d \gg S_f\)), and the kinetics of the reactions are first order

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**Figure 2.** Ligand-receptor system. Left panels \(\tau_{\text{ln}} = 0.01\) (fast stimulation), middle panels \(\tau_{\text{ln}} = 1\), right panels \(\tau_{\text{ln}} = 100\) (slow stimulation). (A–C) Stimulus vs. time \((t_n = t/t_{\text{ref}})\), colors indicate \(x_{\text{max}} = 0.01-0.1-1-10-100-1000\). (D–F) Output \(y\) vs. time \((t_n = t/t_{\text{ref}})\), colors indicate \(x_{\text{max}}\) as before. (G–I) Input-output curves, output \((y)\) vs. input \((x_{\text{max}})\) for the indicated times. (J) EC_{50} vs time for the indicated input dynamics. (K) Hill coefficient \((n_H)\) vs time for the same conditions used in (J).
(i.e., linear) with respect to the concentration of the substrate. In the other extreme regime, \( K_a \) and \( K_d \) are small (\( K_{m,a} \ll S_T \) and \( K_{m,d} \ll S_T \)), the enzymes are saturated by their substrates, and therefore the kinetics of the reactions are zero-order with respect of the concentration of the substrates (only dependent on the amount of enzyme). In this second regime, the steady-state input-output curve is ultrasensitive9, i.e., below a threshold input concentration, changes in input cause almost no change in output. However, past that threshold, a small change in input gives a much larger change in output. Systems with high ultrasensitivity thus behave as molecular switches.

We studied the system in these two extreme regimes (see methods for a detailed description of the parameter scanning), for the case in which the input \( x \) increases fast relative to the timescale of the covalent modification cycle (\( \tau_{Ea,n} \ll 1 \)), and thus it is essentially a step given by \( x(t) = x_{\text{max}} \); and the case in which it is slow compared with the cycle (\( \tau_{Ea,n} \gg 1 \)).

Figure 3(B–M) shows simulations of example parameter sets of each of the four conditions (first- or zero-order, fast or slow input), as well as a parameter scan summarized in Fig. 3N–Q (details in the Supplementary Information). (J–M) EC_{50} (J,L) and \( n_H \) (K,M) vs time for the four examples considered in (B–I). The inset in panel L shows how each curve approaches steady-state (N–Q) EC_{50} (N,P) and \( n_H \) (O,Q) vs time for different parameter sets corresponding to the zero-order regime (N,O) or the first order regime (P,Q). For each scenario (\( K = 100 \) slow, \( K = 100 \) fast, \( K = 0.01 \) slow, \( K = 0.01 \) fast), 100 random parameter sets were scanned, see Supplementary Information for details of the parameter scanning.

Figure 3(B–M) shows simulations of example parameter sets of each of the four conditions (first- or zero-order, fast or slow input), as well as a parameter scan summarized in Fig. 3N–Q (details in the Supplementary Information).
Notably, the input-output curve shifted from right to left over time in all four conditions (Fig. 3F–I). Just as in case of the ligand-receptor system, when the stimulus increased slowly (Fig. 3C,E,G,I), the leftward shift in the input-output curve shifted correspondingly slowly. The above results thus show that CM cycles are shifters. However, we observed interesting differences between CM cycles working in zero-order and first order regimes. The n_H was biphasic in the first-order regime (Fig. 3M and O) with an increase from about 0.85 to 1.35 and a later decrease towards 1. In stark contrast, in the zero-order regime n_H increased monotonically, from about 1 towards its final high steady-state value (~5 to 22 in our simulations) (Fig. 3K and O). This surprising result indicates that the well-known ultrasensitivity of this type of system is not immediately apparent but rather develops over time. In addition, that an ultrasensitive system is a shifter implies that the threshold input required to activate the switch moves over time. Thus, before reaching steady-state, much higher input concentrations would be needed to cross the threshold and flip the switch than at steady-state.

**Coupling shifters (I): a ligand-receptor reaction activating a covalent modification cycle.** To study how the shifting property propagates in a system, we analyzed a toy system composed of the two shifter modules we have described so far: ligand-receptor coupled to a CM cycle (LR-CM) (Fig. 4A). We implemented a full mechanistic model, including enzyme-substrate complexes. (See Supplementary Information for a complete description of the model). We defined two outputs: an upstream output y = (RL + RLS)/R_T, the concentration of ligand-receptor in units of the total amount of receptor, and a downstream output z = S*/S_T, the active substrate in units of the total amount of substrate. We then compared the times at which each output reached steady-state (t_shift and t_shift) and the corresponding EC_{50ss} using different combinations of parameters so that the CM subsystem was either in the first or zero-order regime (see Methods). We made several observations. First, the LR-CM system as a whole was a shifter, since EC_{50ss} shifted from right to left over time (Fig. 4C,D). Second, as expected for a system with no feedback loops (13), EC_{50ss} was always lower than EC_{50ss} (Fig. 4C), indicating that the range of EC_{50ss} spanned by the system was larger than that spanned by the first step alone. This feature is independent of the parameter set used (Fig. 4E). Third, in contrast to the curves of y, the amplitude of the curves of z increased over time. That is, at high input, the output was submaximal if the elapsed time was short enough. Fourth, t_shift was always larger than t_shift (Fig. 4C–F). Interestingly, we found that the ratio between the t_shift for the two outputs correlated well with the ratio of their EC_{50ss} (Fig. 4G): when the EC_{50ss} are nearly identical (a particular case known as Dose-Response Alignment (13,14), t_shift are also identical.

**Coupling shifters (II): a cascade of covalent modification cycles.** In order to evaluate the shifting property in a concrete case, we considered the chain of events involved in oocyte maturation (15). Xenopus oocytes convert a continuously variable stimulus, the concentration of the maturation-inducing hormone progesterone, into an all-or-none biological response: oocyte maturation. The all-or-none character of the response can be accounted for in part by the intrinsic ultrasensitivity of the oocyte’s mitogen-activated protein kinase (MAPK) cascade. In contrast to the zero-order ultrasensitivity we implemented in the CM model above, here the steepness of the input-output curve originates mainly in the distributive nature of the phosphorylation events in the cascade (16).

We used a detailed, mass-action kinetics model of the MAPK cascade previously developed in a combined modeling-experimental study (17) (Fig. 5A). This model describes the dynamics of 22 species participating in 10 reactions. The authors estimated or extracted from cellular and biochemical experiments each of the 37 parameters (See Supplementary Information)).18,19,20 We analyzed the shifter properties of this model. We simulated various input (E_{1tot}) concentrations as steps and monitored the system output (doubly phosphorylated MAPK) over time (Fig. 5B). The output dynamics was sigmoidal, especially evident with low levels of input. That is, there was a long delay until output increased: MAPK-PP began to increase only after 20 minutes with E_{1tot} = 10^{-6} μM. This suggested that transmission of information in this cascade is slow. Remarkably, its input-output curve shifted over time, and did so slowly, only reaching steady-state after about 200 minutes (Fig. 5C–D). As with the case of our LR-CM model above, output was submaximal at short times even in the high input region. Interestingly, the ultrasensitivity of the system was not present at early times, increasing from about 1.4 to around 5 at about 10 minutes (Fig. 5D). This is similar to the behavior displayed by our CM cycle simulated in the zero-order regime (Fig. 3K and O).

In summary, our results indicate that cascades of covalent modification cycles are shifters. Therefore, the range of inputs that these systems would be able to distinguish will strongly depend on the speed of the input they receive (see previous sections) and the relative timing between the cascade and its downstream steps.

**Transcriptional regulation I. A simplified model of gene expression is not a shifter.** Above we have studied how the shifting property propagates through a chain of simple biochemical reactions. However, signaling pathways usually lead to changes in gene expression. Thus, it is also interesting to consider whether a gene expression step could also be a shifter. As a first step, we considered the simplest case of a transcriptional regulatory network where a transcription factor X binds to a promoter region P in the DNA and induces the expression of gene Y into protein Y at a rate γ. Over time, Y is degraded at a rate δ. This process may be described by the following scheme:

\[ P^* \xrightarrow{\gamma} P^* + Y \]  \hspace{1cm} (5.a)

\[ Y \xrightarrow{\delta} \emptyset \]  \hspace{1cm} (5.b)

where P^* stands for the active promoter. This translates into the following differential equation:
$\frac{dY}{dt} = R(X) - \delta Y,$

(6)

where $\gamma = R(X)$ is the regulatory function, $X$ being the transcription factor. Typically $R(X)$ may take the form of a Hill function $aX^n/(X^n + \theta^n)$ with cooperativity $n$ and $EC_{50}$ indicated by $\theta$, depending on the exact way in which it acts.
X binds the DNA. The parameters $a$ and $\delta$ are the maximum synthesis rate of protein Y and the degradation plus dilution parameter, respectively.

In the limit case in which $X$ is not a function of time, considering $R(X) = \frac{aX^n}{(X^n + \theta^n)}$ and that the initial condition is $Y(0) = 0$, equation (6) has a simple solution:

$$Y(t) = \frac{a}{\delta} \frac{X^n}{X^n + \theta^n}[1 - \exp(-\delta t)]$$  

(7)

As has been previously noted, the dynamics of this system only depends on the degradation rate $\delta$, independently of the value of the stimulus $X^n$. $X$ only controls the amplitude of the response $Y$. Not only that, normalizing the input-output curve to the maximum at a given time, results in a time independent input-output curve:

$$\frac{Y(t)}{\frac{a}{\delta} \frac{X^n}{X^n + \theta^n}[1 - \exp(-\delta t)]} = \frac{X^n}{X^n + \theta^n}$$  

(8)

Therefore, this system is not a shifter, and as such may not operate in PRESS mode.

**Transcriptional regulation II. A more detailed model of gene expression is a shifter.** We suspected that the inability of the above transcriptional model to behave as a shifter is not because transcription is intrinsically different from the other biochemical processes, but because the model is oversimplified, omitting one or more reactions crucial for this behavior. An example is the omission of binding between the transcription factor $X$ and the promoter, since, as we explained above, the binding reaction is a shifter. Thus, we expected that a more detailed model of transcriptional regulation would behave as a shifter, at least in some region of the parameter space. To determine if that hypothesis was correct, we considered a model in which a transcription factor $X$ binds a promoter $P$. The bound promoter $P_b$ then transitions to an active promoter $P^*$, which can then synthesize protein $Y$ (Fig. 6A,i,ii).

We described this system by the following equations:

$$\frac{dp}{dt_n} = -x(t_n) \beta_p + \beta_p P_b$$  

(9a)

$$\frac{dp_b}{dt_n} = x(t_n) \beta_p - \beta_p P_b - \theta_{P^*} P_b + \omega_{P^*}$$  

(9b)
\[ \theta \omega = - \delta \delta \]

\[ dp^* \over dt_n = \theta_p \beta_p^* - \omega_p \beta_p^* \tag{9.c} \]

\[ dy \over dt_n = \gamma_p \beta_p^* - Y \tag{9.d} \]
where $x = \frac{X(t)}{\beta \alpha}$ and the conservation law $p + p_b + p^* = 1$, where $p, p_b$ and $p^*$ are the unbound, bound and active promoters, respectively, considering a system with one transcriptional unit (one promoter with one binding site), and the time is expressed in units of a reference time $t_{ref} = 1/\delta$, so that $x = t/t_{ref}$. This reference time is the time-scale of the degradation plus dilution processes. All the rates are in units of the degradation rate $\delta$, resulting in the following parameters $\beta = \beta_b \theta b, \theta = \theta b, \omega = \omega b, \gamma = \gamma b$.

We considered that the concentration of free $X$ is not influenced by the binding and unbinding of transcription factor molecules, so that $X = X_t$. We described the dynamics of transcription factor accumulation by

$$X(t) \approx X(t) = X_{\max}[1 - \exp(-t/\tau)]$$ (10)

where $X_{\max}$ is the maximum value reached by $X$ and $\tau$ is its characteristic time. In this way, and expressing the time in terms of the reference time $t_{ref} = 1/\delta$, we obtained for the dimensionless form of $X, x$:

$$x(t) = \frac{x(t)}{\beta \alpha} = X_{\max}[1 - \exp(-t/\tau_n)]$$ (11)

where $t = t/t_{ref}$ and $\tau_n = \tau/\tau_{ref}$ are the dimensionless time and dimensionless accumulation time of $X$, respectively.

To analyze this model in terms of its kinetic parameters, we first studied its behavior (time course and input-output curve) using an arbitrary parameter set (Fig. 6A.iii, A.iv). Notably, the input-output curve shifted from right to left over time, indicating that this model of transcription does behave as a shifter. As in the cases of the LR-CM model (Fig. 5), the amplitude of the curves changed as well. Thus, to analyze this type of models further, we defined a threshold output (20% of the maximal output possible; this value is arbitrary, it intends to represent a measurable output) and the time at which it was obtained, $t_1$. This time $t_1$ defined the first moment in which there was a reasonable input-output curve. We called the ratio $E_{C50}(t_1)/E_{C50}(t_s)$ of that curve Maximal $E_{C50}$ (ME) (expressed in units of the $E_{C50}$ at steady state). It was also useful to define the Time Window (TW) as the interval between $t_1$ and the time the model reached steady-state ($t_s$):

$$\text{Time Window}(TW) = t_s - t_1$$ (12a)

$$\text{Maximal EC50(ME)} = \frac{E_{C50}(t_1)}{E_{C50}(t_s)}$$ (12b)

A high value of ME indicates that before steady state there is a wide range of distinguishable inputs, and a large TW indicates that there is a considerable time during which pre-steady state information could be influencing the system.

Next, to determine how ME and TW depended on the parameters of the model, we varied two parameters at a time from $10^4$ to $10^8$, while maintaining all the remaining parameters equal to the degradation rate $\delta$ (see Methods for a detailed description of the parameter scan and computation of TW and ME). In this way we generated $10^8$ sets for which we computed ME and TW (Fig. 6A.vi). Several observations are worth mentioning. First, we noted that ME and TW were strongly correlated (Fig. 6A.vi), indicating that the slower the shifter, the larger the distance from the first observable $E_{C50}$ and its steady state value. Second, ME/TW were independent of $\gamma$, the production rate of the protein $Y$, which is constant in this section (Fig. 6E–H). This was expected, since $\gamma$ does not influence the dynamics of promoter activation and only appeared as a proportionality constant in the solution of the ODE for $Y(t)$ (see equations 9a–d). Third, ME and TW increased as either of the two rates that control inactivation of transcription (the dissociation rate of the transcription factor $X$ ($\beta$) or the promoter deactivation rate ($\omega$)) diminished (see for example Fig. 6D). Fourth, the activation ($\theta$) and deactivation ($\omega$) of the promoter were interdependent in the way they controlled ME and TW. When the ratio $K = \omega/\theta$ was small, ME and TW were large (Fig. 6C). Finally, the rate of accumulation of active transcription factor $X$ ($\tau_n$) affected the system as a whole, independently of all the other parameters: if $\tau_n > 1$, ME and TW were large, while if $\tau_n < 1$, then ME and TW were small.

In summary, the above results show that the process of gene expression behaves as a shifter. The shift was larger and it lasted longer if a) the active transcription factor increased slowly, and b) transcription inactivation was slow, both relative to the degradation rate of the gene product. Thus, systems with unstable gene products tend to be better shifters.

Discussion

Previously we have shown that if the input-output curve of a ligand:receptor reaction is measured before it reaches equilibrium, it is shifted to the right of its equilibrium position (Fig. 1C). More specifically, the $E_{C50}$ decreases over time (Fig. 1E), as has been shown experimentally for the yeast $\alpha$-factor:Ste2 pair and for norepinephrine:2A pair and for norepinephrine:2A pair in mammalian cells. In this work, we refer to systems with this property as shifters. Shifting is based on a particular feature of the system: it reaches equilibrium with a characteristic time that is a decreasing function of the input (Fig. 1D). That is, ligands need appreciably more time to reach equilibrium binding at low concentrations than at higher concentrations. Thus, shortly after stimulation, binding is further from equilibrium the lower the ligand concentrations. Consequently, at that early time, while the amount of response may already decrease to steady-state, it is still far less than expected at steady-state at the lowest concentrations. This will result in a steeper saturation curve (i.e. $n_H > 1$; Fig. 1E) and the half-maximal binding ($E_{C50}$) (Fig. 1E) values may also be appreciably higher than expected at a steady-state.

We have also shown that shifting, coupled to a downstream pathway appropriately fast and transient, allows systems to use pre-equilibrium information in order to distinguish stimuli that saturate receptors at equilibrium. We called this ability PRESS (Pre-Equilibrium Sensing and Signaling).
In our previous article\(^1\), the speed of the \textit{shift} depended exclusively on the binding/unbinding rates: the slower the rates, the slower the \textit{shift} and the longer the time window during which the system may exploit pre-equilibrium information. Thus, we speculated that biological systems exposed to step-like stimulus might have evolved slow binding ligand-receptor pairs to extend the range of inputs it can distinguish (the dynamic range). However, in many instances, stimuli increase slowly. Thus, in this work we asked if the requirement for PRESS that the binding reaction be slow may be relaxed by introducing dynamics to the stimulus. Indeed, here we show that input dynamics controls the speed of the \textit{shift}.

The other important goal of the current paper was to find other basic signaling systems that could work as \textit{shifters}, apart from a ligand-receptor binding. We found that covalent modification cycles and simple gene expression systems can \textit{shift} their input-output curve in time. In each case, we studied which conditions improve the \textit{shifting} capacities. The motivation for testing covalent modification cycles as \textit{shifters} comes from the fact that when the enzymes operate in the first order regime, these cycles follow the same mathematical description that is valid for a ligand-receptor system\(^2\). Thus, at least in that limit, a decreasing EC\(_{50}\) versus time was guaranteed. Gene expression systems in which the response time is governed by the degradation/dilution rate only\(^3\), which is usually independent on the input (e.g. transcription factor level) received by this system were not expected to work as \textit{shifters} (see Section Transcriptional regulation I. A simplified model of gene expression is not a \textit{shifter}). However, as soon as the model incorporated additional, more realistic reactions, such as transcription factor binding to the promoter, shifting emerged.

**Ligand-receptor systems.** When receptor binding/unbinding is slow, the EC\(_{50}\) changes correspondingly slowly. Here, we asked how the \textit{shift} in input-output curves is altered when the input has dynamics, particularly when ligand accumulation is slower than the binding/unbinding reaction speed. We considered the case of a ligand that accumulates exponentially. We focused on the impact that the relative time-scales (ligand accumulation versus binding/unbinding) could have on the \textit{shift}. We found that a slow rising input caused a slow \textit{shift}, providing a longer time to a pathway downstream to operate in PRESS mode. Thus, if the stimulus is slow in rising, then the \textit{shift} will also be slow, even in cases where the ligand-receptor interaction is fast. For the case in which ligand accumulation is much slower than the binding/unbinding reaction, we derived a quasi-steady-state approximation leading to an analytical expression for the temporal profile of the EC\(_{50}\). Interestingly, in this limit the slowness of the \textit{shift} is directly controlled by the slowness of ligand accumulation.

**Covalent modification cycle.** In this work we studied CM cycles in their ability to behave as \textit{shifters}, assuming that the activating enzyme increases following an exponential function, for example due to its accumulation after its expression is induced. We considered two extreme regimes: the first-order and the zero-order regimes, both for the cases in which the activating enzyme accumulates fast or slow relative to the timescale of the covalent modification cycle. Mixed scenarios, like saturated kinase and non-saturated phosphatase, and vice versa, are for sure of interest and could lead to different operating regimes\(^3\) but were not covered in the present article.

As obtained for the RL system, we found that when the stimulus increased slowly, the leftward \textit{shift} in the input-output curve was correspondingly slow. We also found that the \textit{shift} was faster in zero-order, indicating that there is less time for PRESS when the enzymes are saturated.

The first-order regime of a CM cycle, mathematically, is similar to the RL system. Thus, we expected that both the EC\(_{50}\) and n\(_{50}\) behaved over time similarly as well. Although the dynamics of the EC\(_{50}\) does indeed behave the same, the dynamics of the n\(_{50}\) was slightly different: it does not decrease monotonically form a 1.42 value towards 1 at steady-state, but it increases from a value close to 1, peaks at 1.42 and only then it declines towards 1.

In the zero-order regime, we found that the ultrasensitivity evidenced in its high n\(_{50}\) takes time to develop, increasing from about 1 towards its final high steady-state value. This surprising result is due to the fact that the system is not actually in the zero-order for backwards reaction at early times, since at these times there is not yet product formed (the substrate for the backwards reaction). Thus, at early times, the system is in first order for one reaction and zero-order for the other.

**Coupling shifters.** We considered the \textit{shifting} properties of two coupled systems, a ligand-receptor coupled to a CM cycle, and a cascade of CM cycles. For the former, we evaluated the \textit{shifting} property throughout the parameter space; for the later, we analyzed a model developed to capture the concrete case of a MAPK cascade involved in oocyte maturation. We found that the LR-CM system as a whole behaves as a \textit{shifter} and also that the \textit{shifting} of the coupled system is slower than that of each step alone. An experimental validation of this result was published some years ago in work studying the activation kinetics of the norepinephrine α2 receptor\(^6\). In this article, the authors measured the activation of this GPCR and of the downstream effector G-protein using a FRET based-method that allows fast measurements. They found that the EC\(_{50}\) of both shifted leftward over time, but that the downstream step exhibited a substantially slower \textit{shift}, in agreement with our Fig. 4D.

Regarding the MAPK cascade, we found a very slow \textit{shift}. As found for a single CM cycle, the ultrasensitivity of a cascade of CM cycles develops over time and thus threshold input required to activate the switch moves over time. In summary, our results indicate that cascades of covalent modification cycles are \textit{shifters}. Therefore, the range of inputs that these systems would be able to distinguish will strongly depend on the speed of the input they receive (see previous sections) and of the relative timing between the cascade and its downstream outputs steps.

**Transcriptional regulation.** We found that this system is not a \textit{shifter} only in the limit case in which its dynamics simply depends on the degradation rate. The reason for this is that this limit case omits reactions that are crucial for the shifting behavior, like the binding between the transcription factor and the promoter. A more detailed model for transcriptional regulation did behave as a \textit{shifter}. The \textit{shift} was larger and it lasted longer if a)
the active transcription factor increased slowly, and b) transcription inactivation was slow, all this relative to the degradation rate of the gene product. Thus, systems with unstable gene products tend to be better shifters.

Pre-steady-state conditions in signaling. Others have also payed attention to pre-steady-state conditions in input-output curves. For example, Charlton and Vauquelin discussed the potential issues associated with the interpretation of receptor pharmacology using calcium assays. Particularly, they focused on the influence of non-equilibrium conditions in that rapid signaling system, observing that the ligand pharmacology at early times can be different from that in equilibrium. Also, Nyman et al. analyzed the measured input-output curve for the \( \beta \)-adrenergic system in the heart, where they were able to re-interpret and thus correct existing data regarding drug potency by using kinetics simulations. Thus, the theoretical framework we developed in this paper allows for reinterpretation of input-output curves when some or all of the data points do not correspond to steady-state measurements.

Importantly for some signaling systems, pre-steady-state considerations may provide a temporal window to use information that is no longer available at steady-state. The PRESS mechanism is an example of this situation. Similarly, Cattoni et al. studied how to discriminate between negative cooperativity and ligand binding to multiple independent sites, two different molecular explanations for the same experimental results when obtained with steady-state experiments. They found that the two ligand binding mechanisms can be readily distinguished by a pre-equilibrium analysis.

Shifting mechanism might be widespread. Summarizing, under physiological conditions, receptor binding and activation, and subsequent signaling are often dependent on very short stimuli and usually do not reach a steady-state. Therefore, it is relevant to study receptor activation and signaling, its potencies, efficacies and sensitivities under non-equilibrium conditions.

Our results here apply particularly to systems that need to adapt their input dynamic range (the range of inputs for which the system is able to generate distinguishable outputs) according to different scenarios. Binding/unbinding reactions, covalent modification cycles and gene expression regulation are the building blocks of cellular regulation, underlying all processes in the cell, from basic homeostatic maintenance of cellular structures and energy balance, to cell division, cell differentiation. The fact that we have shown that these elements all may behave as shifters indicates that attention to dynamics of each of these processes might shed interesting insights into fundamental biological processes.

Methods Numerical integration of ODEs. ODEs were integrated using the ode23s function from Matlab.

Covalent-Modification cycle: computation of the Hill coefficient and parameter scanning. We used the definition \( n_H = \log(81)/\log(EC_{90}/EC_{10}) \). To calculate \( n_H \) at a given time \( t \), we calculated \( EC_{90} \) and \( EC_{10} \) at time \( t \). Hence, being \( y(x, t = t_0) \) the Input(x)-Output(y) curve for a fixed time \( t_0 \), we computed \( EC_{90}(t = t_0) \) and \( EC_{10}(t = t_0) \) from: \( y(x = EC_{90}, t = t_0) = 0.9 \max \ y(x, t = t_0) \) and \( y(x = EC_{10}, t = t_0) = 0.1 \max \ y(x, t = t_0) \) respectively.

Therefore, the Hill coefficient at a given time \( t_0 \) is computed as: \( n_H(t = t_0) = \log_{10}(81)/\log_{10}(EC_{90}(t = t_0)/EC_{10}(t = t_0)) \). Input-Output curves can be obtained analytically only in some cases, such as the ligand-receptor system (equations S3, S6 and S7). Otherwise, they must be computed numerically.

Therefore, to compute \( EC_{90} \) and \( EC_{10} \) at \( t_0 \), we used the discrete grid of pairs \((x_i, y_i)\) and connected them by linear interpolation. The same method was used in other parts of the paper that need Hill coefficient calculation.

Parameter Scanning. We scanned parameter values randomly with log uniform distribution within the intervals defined on Table S1. We used only random parameter sets whose \( K_s \) and \( K_d \) fell within the intervals.

\[
0.009 < K_p, K_d < 0.011, \quad \text{for the zero-order set}
\]

\[
90 < K_p, K_d < 110, \quad \text{for the first-order set}
\]

\( K_s \) and \( K_d \) correspond to the Michaelis-Menten constants \( K_s = (d_1 + k_d)/aS_p, K_d = (d_1 + k_d)/aS_p \), in units of the total amount of substrate. We kept the first 100 sets satisfying the mentioned conditions. We then simulated the CM cycle using each of those sets and obtained input-output curves for 50 different times within the interval \([10^{-3}, 10^3]\), for fast and slow input.

A ligand-receptor activates a covalent modification cycle: parameter scanning. To obtain different parameter sample sets, we used the same criterion as with the CM cycles. Time was in units of \( k_{ad} \) (i.e. parameters are in units of the unbinding rate) and the input in units of the dissociation constant of the binding reaction. We scanned \( a_1, a_2, d_1, d_2, k_1, k_2, S_T \) and \( E_T \) randomly, from a uniform distribution within the intervals defined in Table S1, and \( R_0 \) in the interval \([0,1500]\). This interval was chosen based on the total amount of substrate (see Table S1) to cover three possible scenarios: \( R_T \ll S_T, R_T \sim S_T, R_T \gg S_T \). Under these considerations, we kept two groups of parameter sets that satisfied

\[
0.0099 < K_{a,d} < 0.0101, \quad \text{for the zero-order set}
\]

\[
99 < K_{a,d} < 101, \quad \text{for the first-order set}
\]

In Fig. 4 we show the behavior of 200 parameter sets for each group that produced an output \( S' \) with an amplitude greater than 0.1 (10% of \( S_T \)).
Transcriptional regulation. Parameter scanning. To quantify the effect that each parameter has on the
shifter, we scanned parameters $\beta_{\gamma}$, $\theta$, $\omega_{\gamma}$, $\gamma_{\theta}$, and $\sigma$ with the following approach. We varied two parameters in a grid of $150 \times 150$ in log scale from $10^{-2}$ to $10^5$ while maintaining the remaining parameters fixed to unity (i.e., kinetic rates and inverse time scale $\sigma$ were equal to the degradation rate $\delta$). For each point on the grid, we computed the Maximal EC$_{50}$ and the Time Window, as defined in the main text. We studied all possible pair of parameters.

Computing $t_1$ and $t_{\nu}$. To compute the first observable time, we defined $t_1$ as the time that satisfies

$$y(x = \infty, t = t_1) = 0.2y(x = \infty, t = \infty)$$

the time at which the (dimensionless) output reached 20% of the maximal output. We obtained this value $t_1$ using a time grid of $N_{t_1} = 1000$ and varying the time in linear scale from 0 to 50 (time in dimensionless units).

Having $t_1$, we computed dose-response curves for different times. We varied the dose in log scale from $10^{-2}$ to $10^4$ with 50 values, and time in log scale from $t_1$ to $10^5$ using 500 values. Such a density in the time grid is needed to accurately compute EC$_{50}$ versus time.

We obtained $t_{\nu}$ as the first time at which $\log EC_{50}(t) < \log EC_{50}(1) < 10^{-2}$. As $EC_{50}(t)$ decreases monotonically with an asymptotic value, its log derivative is negative and goes asymptotically to zero. Hence, we considered that the dose response curve reached steady state when $\frac{d\log EC_{50}(t)}{d\log t} < 10^{-2}$. Note that, strictly, steady state is reached at different times for different inputs. However, for practical reasons, we used the time that the $EC_{50}$ reached steady state as a proxy for the whole input-output curve. Having $EC_{50}(t_1)$, $t_1$, and $t_{\nu}$ we used equation 12 to compute Maximal EC$_{50}$ and Time Window.

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Author Contributions
A.C.V. and A.C.L. designed the project. J.D.B. and A.C.V. performed all mathematical analysis and simulations. J.D.B., A.C.V. and A.C.L. analyzed the results and prepared the manuscript.

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