Measurement of the Semileptonic Decays $B \to D\tau^-\nu_\tau$ and $B \to D^{*}\tau^-\nu_\tau$

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We present measurements of the semileptonic decays $B^- \to D^0 \tau^- \nu_\tau$, $B^- \to D^{*0} \tau^- \nu_\tau$, $\bar{B}^0 \to D^{*+} \tau^- \nu_\tau$, and $\bar{B}^0 \to D^{*0} \tau^- \nu_\tau$, which are sensitive to non-Standard Model amplitudes in certain scenarios. The data sample consists of $232 \times 10^6$ $T(4S) \to B\bar{B}$ decays collected with the BABAR detector at the PEP-II $e^+e^-$ collider. We select events with a $D$ or $D^*$ meson and a light lepton ($\ell = e$ or $\mu$) recoiling against a fully reconstructed $B$ meson. We perform a fit to the joint distribution of lepton momentum and missing mass squared to distinguish signal events from the backgrounds, predominantly $B \to D^{(*)0} \ell^- \nu_\ell$. We measure the branching-fraction ratios $R(D) \equiv B(B \to D\tau^- \nu_\tau)/B(B \to D^*\ell^- \nu_\ell)$ and $R(D^*) \equiv B(B \to D^*\tau^- \nu_\tau)/B(B \to D^{*0}\ell^- \nu_\ell)$ and, from a combined fit to $B^-$ and $\bar{B}^0$ channels, obtain the results $R(D) = (41.6 \pm 11.7 \pm 5.2)\%$ and $R(D^*) = (29.7 \pm 5.6 \pm 1.8)\%$, where the uncertainties are statistical and systematic. Normalizing to measured $B^- \to D^{(*)0} \ell^- \nu_\ell$ branching fractions, we obtain $B(B \to D\tau^- \nu_\tau) = (0.86 \pm 0.24 \pm 0.11 \pm 0.06)\%$ and $B(B \to D^*\tau^- \nu_\tau) = (1.62 \pm 0.31 \pm 0.10 \pm 0.05)\%$, where the additional third uncertainty is from the normalization mode. We also present, for the first time, distributions of the lepton momentum, $|p_\ell|$, and the squared momentum transfer, $q^2$.

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I. INTRODUCTION

Semileptonic decays of $B$ mesons to the $\tau$ lepton—the heaviest of the three charged leptons—provide a new source of information on Standard Model (SM) processes [1, 2, 3], as well as a new window on physics beyond the SM [4, 5, 6, 7, 8, 9]. In the SM, semileptonic decays occur at tree level and are mediated by the $W$ boson, but the large mass of the $\tau$ lepton provides sensitivity to additional amplitudes, such as those mediated by a charged Higgs boson. Experimentally, $b \to c\tau^-\overline{\nu}_\tau$ decays are challenging to study because the final state contains not just one, but two or three neutrinos as a result of the $\tau$ decay.

Theoretical predictions for semileptonic decays to exclusive final states require knowledge of the form factors, which parameterize the hadronic current as functions of $q^2 = |p_B - p_{\nu_L}|^2$. For light leptons $\ell = e, \mu$, there is an effectively one form factor for $B \to D\ell^-\overline{\nu}_\ell$, while there are three for $B \to D^*\ell^-\overline{\nu}_\ell$. If a $\tau$ lepton is produced instead, one additional form factor enters in each mode. The form factors for $B \to D^{(*)}\ell^-\overline{\nu}_\ell$ decays involving light leptons have been measured [10, 11, 12], providing direct information on four of the six form factors. Heavy quark symmetry (HQS) relations [13] allow one to express the two additional form factors for $B \to D^{(*)}\tau^-\overline{\nu}_\tau$ in terms of the form factors measurable from decays with the light leptons. With sufficient data, one could probe the additional form factors and test the HQS relations.

Branching fractions for semileptonic $B$ decays to $\tau$ leptons are predicted to be smaller than those to light leptons. Calculations based on the SM predict $B(B^0 \to D^+\tau^-\overline{\nu}_\tau) = (0.69 \pm 0.04)\%$ and $B(B^0 \to D^{*+}\tau^-\overline{\nu}_\tau) = (1.41 \pm 0.07)\%$ [8], which account for most of the predicted inclusive rate $B(B \to X_c\tau^-\overline{\nu}_\tau) = (2.30 \pm 0.25)\%$ [2] (here, $X_c$ represents all hadronic final states from the $b \to c$ transition). In multi-Higgs doublet models [4, 5, 6, 7, 8], substantial departures, either positive or negative, from the SM decay rate could occur for $B(B \to D\tau^-\overline{\nu}_\tau)$, while smaller departures are expected for $B(B \to D^*\tau^-\overline{\nu}_\tau)$. Thus, measurements of $B(B \to D\tau^-\overline{\nu}_\tau)$ are more sensitive to non-SM contributions than either $B(B \to D^*\tau^-\overline{\nu}_\tau)$ or the inclusive rate. In addition to the branching fractions, several other observables are sensitive to possible non-SM contributions, including $q^2$ distributions and $D^*$ and $\tau$ polarization [1, 5, 6, 8, 14].

The first measurements of semileptonic $b$-hadron decays to $\tau$ leptons were performed by the LEP experiments [15] operating at the $Z^0$ resonance, yielding an average [16] inclusive branching fraction $B(b_{\text{had}} \to X\tau^-\overline{\nu}_\tau) = (2.48 \pm 0.26)\%$, where $b_{\text{had}}$ represents the mixture of $b$-hadrons produced in $Z^0 \to b\overline{b}$ decays. The Belle experiment has reported $B(B^0 \to D^{*+}\tau^-\overline{\nu}_\tau) = (2.02^{+0.40}_{-0.37} \pm 0.37)\%$ [17].

The BABAR Collaboration has presented a measurement of the branching fractions for $B \to D\tau^-\overline{\nu}_\tau$ and $B \to D^*\tau^-\overline{\nu}_\tau$ for both charged and neutral $B$ mesons [18]. In this article, we describe the analysis in greater detail, with particular emphasis on several novel features of the event selection and fit technique. We also present distributions of two important kinematic variables, the lepton momentum, $|p_\ell^{|}$, and the squared momentum transfer, $q^2$.

A. Analysis overview and strategy

We determine the branching fractions of four exclusive decay modes: $B^\pm \to D^{0}\tau^-\overline{\nu}_\tau$, $B^\pm \to D^{*+}\tau^-\overline{\nu}_\tau$, $B^0 \to D^+\tau^-\overline{\nu}_\tau$, and $B^0 \to D^{*+}\tau^-\overline{\nu}_\tau$, each of which is measured as a branching-fraction ratio $R$ relative to the corresponding $e$ and $\mu$ modes. To reconstruct the $\tau$, we use the decays $\tau^- \to e^-\overline{\nu}_e\nu_e$ and $\tau^- \to \mu^-\overline{\nu}_\mu\nu_e$, which are experimentally the most accessible. The main challenge of the measurement is to distinguish $B \to D^{(*)}\tau^-\overline{\nu}_\tau$ decays, which have three neutrinos, from $B \to D^{(*)}\ell^-\overline{\nu}_\ell$ decays, which have the same observable final-state particles but only one neutrino.

The analysis strategy is to reconstruct the decays of both $B$ mesons in the $Y(4S) \to B\overline{B}$ event, providing powerful constraints on unobserved particles. One $B$ meson, denoted $B_{\text{tag}}$, is fully reconstructed in a purely hadronic decay chain. The remaining charged particles and photons are required to be consistent with the products of a $b \to c$ semileptonic $B$ decay: the daughter charm meson (either a $D$ or $D^*$) and a lepton ($e$ or $\mu$). The lepton may be either primary or from $\tau^- \to \ell^-\overline{\nu}_\ell\nu_e$. To distinguish signal events from the normalization modes $B \to D^{(*)}\ell^-\overline{\nu}_\ell$, we calculate the missing four-momentum,

$$p_{\text{miss}} = p_{\ell^+e^-} - p_{\text{tag}} - p_{D^{(*)}\ell} - p_\ell$$

(1)
of any particles recoiling against the observed $B_{\text{tag}} + D^{(*)}\ell$ system. A large peak at zero in $m^2_{\text{miss}} = p_{\text{miss}}^2$ corresponds to semileptonic decays with one neutrino, whereas signal events produce a broad tail out to $m^2_{\text{miss}} \sim 8$ (GeV/c$^2$)$^2$.

To separate signal and background events, we perform a fit (described in Section VII) to the joint distribution of $m^2_{\text{miss}}$ and the lepton momentum ($|p_\ell^{|}$) in the rest frame of the $B$ meson. In signal events, the observed lepton is the daughter of the $\tau$ and typically has a soft spectrum; for most background events, this lepton typically has higher momentum. The fit is performed simultaneously in eight channels, with a set of constraints relating the event yields between the channels. The fit is designed to maximize the sensitivity to the $B \to D\tau^-\overline{\nu}_\tau$ signals by using events in the $D^*\ell^-\overline{\nu}_\ell$ channels to constrain

1 Charge-conjugate modes are implied throughout.
2 Throughout this article, we use the symbol $\ell$ to refer only to the light charged leptons $e$ and $\mu$.
3 The symbol $D^{(*)}$ refers either to a $D$ or a $D^*$ meson.
the dominant backgrounds, \( B \rightarrow D^* \tau^- \nue \), feed-down, in which the final-state \( D^* \) meson is not completely reconstructed. Similarly, we use a set of \( D^{**} \) control samples to constrain the feed-down background to both the \( D^\tau^- \nue \) and \( D^* \tau^- \nue \) signals.\(^4\)

We perform a relative measurement, extracting both signal \( B \rightarrow D^{(*)} \ell^- \nue \) and normalization \( B \rightarrow D^{(*)} \ell^- \nu \) yields from the fit to obtain the four branching-fraction ratios \( R(D^0), R(D^+), \) and \( R(D^{(*)}_s) \), where, for example, \( R(D^{(*)}_s) \equiv B(B^- \rightarrow D^{(*)}_s \tau^-)/B(B^- \rightarrow D^{(*)}_s \nu) \). In the ratio, many systematic uncertainties cancel, either partially or completely. These ratios are normalized such that \( \ell \) represents only one of \( e \) or \( \mu \); however, both light lepton species are included in the measurement. We multiply these branching-fraction ratios by previous measurements of \( B(\ell \rightarrow D^{(*)}) \) to derive absolute branching fractions.

II. THE BABAR DETECTOR AND DATA SETS

We analyze data collected with the BABAR detector at the PEP-II \( e^+e^- \) storage rings at the Stanford Linear Accelerator Center. PEP-II is an asymmetric-energy \( B \) factory, colliding 9.0 GeV \( e^- \) with 3.1 GeV \( e^+ \) at a center-of-mass energy of 10.58 GeV, corresponding to the \( \Upsilon(4S) \) resonance. The data sample used consists of 208.9 fb\(^{-1} \) of integrated luminosity recorded on the \( \Upsilon(4S) \) resonance between 1999 and 2004, yielding \( 232 \times 10^6 \Upsilon(4S) \rightarrow \BB \) decays. This data sample can be divided into two major periods: Runs 1–3, comprising 109.0 fb\(^{-1} \) taken from 1999 to June 2003, and Run 4, comprising 99.9 fb\(^{-1} \) taken from September 2003 to July 2004. The accelerator background conditions were significantly different between Runs 1–3 and Run 4, which could affect missing-energy analyses such as this one; for this reason, the two running periods have been independently validated, and the fraction of signal-like events found in the Run 4 sample is used as a crosscheck of the results, as described in Section III.

The BABAR detector is a large, general-purpose magnetic spectrometer and is described in detail elsewhere.\(^10\) Charged particle trajectories are measured in a tracking system consisting of a five-layer double-sided silicon strip detector and a 40-layer drift chamber, both of which operate in the 1.5 T magnetic field of a superconducting solenoid. A detector of internally reflected Cherenkov light (DIRC) is used to measure charged particle velocity for particle identification (PID). An electromagnetic calorimeter (EMC), consisting of 6580 CsI(Tl) crystals, is used to reconstruct photons and in electron identification. The steel flux return of the solenoid is segmented and instrumented with resistive plate chambers (IFR) for muon and neutral hadron identification.

All detector systems contribute to charged particle identification. Ionization energy losses in the tracking systems and the Cherenkov light signature in the DIRC are used for all charged particle types. Electrons are also identified on the basis of shower shape in the EMC and the ratio of energy deposited in the EMC to the track momentum. Muon identification is based on a minimum-ionization energy deposit in the EMC and on the measured interaction length in the IFR.

This analysis relies on measurement of the missing momentum carried off by multiple neutrinos, and the large solid angle coverage (hermeticity) of the detector is therefore crucial. The tracking system, calorimeter, and IFR cover the full azimuthal range and the polar angle range from approximately 0.3 \( < \theta < 2.7 \) rad in the laboratory frame, corresponding to a \( \Upsilon(4S) \) center-of-mass coverage of approximately 90% (the direction \( \theta = 0 \) corresponds to the direction of the high-energy beam, and therefore to the \( \Upsilon(4S) \) boost). The DIRC fiducial volume is slightly smaller, corresponding to a center-of-mass frame coverage of about 84%.

Within the active detector volume, the efficiency for reconstructing charged tracks and photons is very high, typically greater than 95% over most of the momentum range. At low momenta, however, the reconstruction efficiency drops off, leading to an increased contribution from feed-down processes to which special attention is paid throughout this analysis. Feed-down occurs when the photon from \( D^* \rightarrow D \gamma \) or the \( \pi^0 \) from \( D^* \rightarrow D \pi^0 \) is not reconstructed (in the case of the \( \pi^0 \), either one or both of the photons from \( \pi^0 \rightarrow \gamma \gamma \) may be missed). Care must therefore be taken to avoid confusing \( D^* \) feed-down events for \( D \) signals.

We use a Monte Carlo simulation (MC) of the production and decay of signal and background events based on EvtGen\(^20\). A sample of simulated inclusive \( \BB \) events equivalent to about five times the integrated luminosity is used to study backgrounds and to optimize event selection criteria. Large samples of many individual semileptonic \( B \) decays (discussed in Section III) are used to parameterize the distributions of variables used in the fit. Final-state radiation is simulated using PHOTOS\(^21\). Simulation of the detector response is performed with GEANT\(^22\) and the resulting efficiencies and resolutions are validated in multiple data control samples.

III. SEMILEPTONIC DECAY MODELS

In the SM, the matrix element for a semileptonic \( B \) meson decay can be written as

\[
\mathcal{M} \left( B \rightarrow D^{(*)} (\ell^-/\tau^-) \nu \right) = -i \frac{g^2}{8m_W^2} V_{cb} L^\mu H_\mu \ , \quad (2)
\]

where \( g \) is the weak coupling constant, \( m_W \) the \( W \) mass,

\(^4\) Through this paper, we use the symbol \( D^{**} \) to represent all charm resonances heavier than the \( D^* \) (2010), as well as non-resonant \( D^{(*)}_n \pi \) systems with \( n \geq 1 \).
$V_{cb}$ the quark mixing matrix element, and $L^\mu$ and $H_\mu$ are the leptonic and hadronic currents, respectively. Here, we have used a simplified form for the $W$ propagator appropriate for energies less than $m_W$. The leptonic current is exactly known,

$$L^\mu = \overline{\nu} \gamma^\mu (1 - \gamma_5) \nu ,$$  
and the hadronic current is given by

$$H_\mu = \langle D^{(*)} | \overline{q} \gamma_\mu (1 - \gamma_5) q | B \rangle .$$

In the case of a $B \to D$ transition, the axial-vector part of the current does not contribute to the decay, and we may write the hadronic current in terms of two form factors $f_+ (q^2)$ and $f_- (q^2)$:

$$\langle D | V^\mu | B \rangle = (p + p')^\mu f_+ (q^2) + (p - p')^\mu f_- (q^2) ,$$

with $V^\mu = \overline{\tau} \gamma^\mu b$ and where $p$ and $p'$ are the four-momenta of the $B$ and $D$ mesons, respectively. For the $B \to D^*$ transition, the axial-vector term contributes to the decay as well, and we write the hadronic current in terms of form factors $V (q^2)$, $A_1 (q^2)$, $A_2 (q^2)$, $A_3 (q^2)$, and $A_0 (q^2)$:

$$\langle D^* | V^\mu - A^\mu | B \rangle = \frac{2 i e m_{\nu \alpha \beta}}{m_B + m_{D^*}} \varepsilon^{\mu \nu \alpha \beta} p_\alpha p_\beta V (q^2) -$$

$$(m_B + m_{D^*}) \varepsilon^\mu A_1 (q^2) + \frac{\varepsilon^\mu q}{m_B + m_{D^*}} (p + p')^\mu A_2 (q^2) +$$

$$2 m_{D^*} \varepsilon^\mu q A_3 (q^2) - 2 m_{D^*} \varepsilon^\mu q A_0 (q^2) ,$$

where $A^\mu = \overline{\tau} \gamma^\mu b$ and $\varepsilon$ is the $D^*$ polarization vector. The form factor $A_3 (q^2)$ is related to two other form factors as

$$A_3 (q^2) = \frac{m_B + m_{D^*}}{2 m_{D^*}} A_1 (q^2) - \frac{m_B - m_{D^*}}{2 m_{D^*}} A_2 (q^2)$$

so that there are only four independent form factors.

In the limit of massless leptons, any terms proportional to $q^\mu \equiv (p - p')^\mu$ vanish when the hadronic current is contracted with the leptonic current. For this reason, the contributions from the form factors $f_- (q^2)$ and $A_0 (q^2)$ are essentially negligible for electrons and muons, as mentioned above.

Semileptonic decays are simulated using the ISGW2 model \textsuperscript{23}, except for $B \to D^* \ell^- \overline{\nu}_\ell$ decays, which use an HQET model with a linear form factor expansion \textsuperscript{24}, and nonresonant $B \to D^{(*)} \pi^- \nu_{\pi}$ decays, which use the model of Goity and Roberts \textsuperscript{25}. We reweight both signal $B \to D^{(*)} \tau^- \overline{\nu}_\tau$ and normalization $B \to D^{(*)} \ell^- \overline{\nu}_\ell$ events \textsuperscript{26} so that the decay distributions follow the Caprini-Lellouch-Neubert (CLN) form factor model \textsuperscript{27} with parameters measured in data. We use $\rho_2^\pm = 1.17 \pm 0.18$ \textsuperscript{28} for $B \to D^0 \overline{\nu}_\ell$ and $B \to D \tau^- \overline{\nu}_\tau$ decays, and we use $R_1 = 1.417 \pm 0.061 \pm 0.044$, $R_2 = 0.836 \pm 0.037 \pm 0.022$, and $\rho_3 A_1 = 1.179 \pm 0.048 \pm 0.028$ \textsuperscript{21} for $B \to D^* \ell^- \overline{\nu}_\ell$ and $B \to D^* \tau^- \overline{\nu}_\tau$ decays.\textsuperscript{5} Variation of these form factors is taken into account as a systematic uncertainty, including the correlations between the three $B \to D^*$ form factor parameters.

Figures \textsuperscript{11-13} show distributions of three kinematic variables important to this analysis, all generated using the CLN form factor parameterization with parameters given above. Figure \textsuperscript{11} compares $q^2$ distributions between the signal and normalization modes. Signal events must satisfy $q^2 > m_\ell^2$, leading to qualitatively different $q^2$ spectra for signal and normalization events; this feature is exploited in the event selection and in validation studies. Figure \textsuperscript{12} shows distributions of lepton energy in the $B$ meson rest frame. While the $\tau^-$ lepton in signal events typically has high energy (due to its mass), the secondary lepton $\ell^-$ typically has much lower energy than either the $\tau^-$ or the primary lepton in $B \to D^{(*)} \ell^- \overline{\nu}_\ell$ events. This low lepton energy leads to a lower reconstruction efficiency for signal leptons than those in the normalization modes. Figure \textsuperscript{13} shows distributions of $m^2_{\text{miss}}$ for the two signal modes, which, due to the three neutrinos in these events, forms a broad structure up to very large $m^2_{\text{miss}}$.

IV. EVENT RECONSTRUCTION AND SELECTION

All event selection requirements (as well as the fit procedure described in Section VII) are defined using simulated events or using control samples in data that exclude the signal region in order to avoid any potential sources of bias. About 60% of the $B\bar{B}$ MC sample is used in optimizing the event selection, while the remaining 40% is used as an independent validation of the selection and fitting procedures.

Most of the selection criteria described here are optimized to maximize the quantity $S/\sqrt{S + B}$, where $S$ and $B$ are the expected signal and background yields in the large $m^2_{\text{miss}}$ region of our data sample, assuming Standard Model branching fractions for signal decays. The requirement on $\Delta E$ of the $B_{\text{tag}}$ candidate (defined below) was initially optimized in the same way, but was tightened because fits to MC samples indicated that events at large $|\Delta E|$ contributed to biases in the signal extraction. The final selection corresponds to a compromise

\footnotetext[5]{The parameters $R_1$ and $R_2$ are not included in the model of Caprini, Lellouch, and Neubert \textsuperscript{23} to model the $B \to D^*$ form factors, we adopt the formalism used in \textsuperscript{12}, Eqs. (13–14), where the leading terms in these form factor ratio expansions are taken as free parameters. We use independent slope parameters $\rho_2^\pm$ and $\rho_3 A_1$ for the $B \to D$ and $B \to D^*$ form factors, respectively, treating the two sets of form factors as uncorrelated.}
between the statistical $S/\sqrt{S+B}$ optimization and the systematic effects due to this bias.

### A. $B_{\text{tag}}$ Reconstruction

We reconstruct $B_{\text{tag}}$ candidates in 1114 final states $B_{\text{tag}} \to D(*)Y^\pm$ with an algorithm that has been used previously at BaBar for a number of analyses, especially those dependent on measuring missing momentum [29]. These final states arise from the large number of ways to reconstruct the $D$ and $D^*$ mesons within the $B_{\text{tag}}$ candidate and the possible pion and kaon combinations within the $Y^\pm$ system. Tag-side $D$ candidates are reconstructed as $D_{\text{tag}}^0 \to K^-\pi^+, K^-\pi^+\pi^0, K^-\pi^+\pi^+\pi^-$, and $K^0_s\pi^+$, and as $D_{\text{tag}}^+ \to K^-\pi^+\pi^0, K^-\pi^+\pi^+\pi^-$, $K^0_s\pi^+$, $K^0_s\pi^+\pi^-\pi^+$, and $K^0_s\pi^+\pi^-\pi^0$. Tag-side $D^*$ candidates are reconstructed as $D_{\text{tag}}^* \to D_{\text{tag}}^0 \pi^0$ and $D_{\text{tag}}^0 \gamma$ and as $D_{\text{tag}}^{*+} \to D_{\text{tag}}^0 \pi^+$. The $Y^\pm$ system may consist of up to six light hadrons ($\pi^\pm$, $\pi^0$, $K^\pm$, or $K^0_s$). In both the $D_{\text{tag}}^{(*)}$ and $Y^\pm$ systems, we reconstruct $\pi^0 \to \gamma\gamma$ and $K^0_s \to \pi^+\pi^-$. The terms “loose” and “tight” refer to the relative signal-to-background discrimination of various PID criteria. Loose criteria are chosen to have high efficiency, and have relatively high background rates as well; tight criteria have lower background but also higher signal efficiency. The optimal choice of criteria depends on the particle type and on the a priori purity of the sample, and is therefore different for each reconstruction channel.

$\Delta m_{\text{ES}} = \sqrt{s/4 - |p_{\text{tag}}|^2}$ (8)
and

$$\Delta E = E_{\text{tag}} - \sqrt{s}/2 ,$$

(9)

where $\sqrt{s}$ is the total $e^+e^-$ energy, $|p_{\text{tag}}|$ is the magnitude of the $B_{\text{tag}}$ momentum, and $E_{\text{tag}}$ is the $B_{\text{tag}}$ energy, all defined in the $e^+e^-$ center-of-mass frame. For correctly reconstructed $B_{\text{tag}}$ candidates, $m_{\text{ES}}$ is equal to the $B$ meson mass, with a resolution of about 2.5 MeV/c$^2$, and $\Delta E$ is equal to zero, with a resolution of about 18 MeV.

For each $D_{\text{tag}}^{(*)}$ “seed” candidate, we use a recursive algorithm to identify candidate $Y^\pm$ systems. Light hadrons from the remaining tracks and photons in the event are added to the $Y^\pm$ system, one at a time. If the resulting values of $m_{\text{ES}}$ and $\Delta E$ for the $D_{\text{tag}}^{(*)}Y^\pm$ candidate are close to the nominal values, the $B_{\text{tag}}$ candidate is accepted. If the value of $\Delta E$ is too large, the light hadron just added is removed from the $Y^\pm$ system, since continuing to add particles to this $Y^\pm$ candidate will increase $\Delta E$ further. The algorithm then continues recursively with the remaining particles in the event, adding and removing light hadrons to the $Y^\pm$ system according to $m_{\text{ES}}$, $\Delta E$, and the $Y^\pm$ system topology. This algorithm is semiexclusive, meaning that particles in the $Y^\pm$ system are not constrained to intermediate resonance states. Because of this, the yield is significantly higher than exclusive $B$ reconstruction, while the purity is somewhat lower. In this analysis, however, since we exclusively reconstruct the second $B$ meson in the event, the purity of our final sample is substantially improved with respect to the raw $B_{\text{tag}}$ sample.

We require $m_{\text{ES}} > 5.27$ GeV/c$^2$ and $|\Delta E| < 72$ MeV, corresponding to $\pm 4\sigma$ in $\Delta E$ and $-4\sigma$ in $m_{\text{ES}}$ (the kinematic limit $m_{\text{ES}} < \sqrt{s}/2$ provides an effective $+4\sigma$ requirement). We reconstruct $B_{\text{tag}}$ candidates with an efficiency of 0.2% to 0.3%. Figure 4 shows distributions of $m_{\text{ES}}$ for selected $B_{\text{tag}}$ candidates both before and after the signal-side reconstruction. We make no attempt at this stage to select a single $B_{\text{tag}}$ among multiple reconstructed candidates: this decision is made after reconstructing the signal side as well.
B. Reconstruction of the Signal \( B \)

For the \( B \) meson decaying semileptonically, we reconstruct \( D^{(*)} \) candidates in the modes \( D^0 \to K^-\pi^+ \), \( K^-\pi^+\pi^- \), \( K^-\pi^+\pi^-\pi^- \), \( K^0\pi^+\pi^- \), \( D^+ \to K^-\pi^+\pi^+ \), \( K^-\pi^+\pi^-\pi^0 \), \( K^0\pi^+\pi^- \), \( K^-\pi^+\pi^- \), \( D^0 \to D^0\pi^0 \), \( D^0\pi^0 \); and \( D^{*+} \to D^0\pi^0\pi^+ \), \( D^+\pi^0 \). We reconstruct \( K^0_s \) mesons as \( K^0_s \to \pi^+\pi^- \) with \( 491 < m_{\pi^+\pi^-} < 506 \) MeV/c\(^2\), corresponding to \( \pm 3\sigma \). We reconstruct \( \pi^0 \) mesons as \( \pi^0 \to \gamma\gamma \), requiring \( 90 < m_{\gamma\gamma} < 165 \) MeV/c\(^2\) for the soft \( \pi^0 \) used to reconstruct \( D^* \to D\pi^0_{soft} \), and requiring \( E_{\gamma\gamma} > 200 \) MeV and \( 125 < m_{\gamma\gamma} < 145 \) MeV/c\(^2\) for a \( \pi^0 \) used to reconstruct a \( D \) meson; the mass intervals correspond to \( \pm 3\sigma \) in both cases and are different because the resolution is poorer at low energies. Charged kaon candidates are required to satisfy tight PID criteria with a typical efficiency of 85% while rejecting 98% of pions. Charged pion candidates are required to satisfy loose PID criteria with a typical efficiency of 97% while rejecting 88% of kaons. \( D \) (\( D^* \)) candidates are selected within \( 4\sigma \) of the \( D \) mass (\( D^* - D \) mass difference); as on the tag side, \( \sigma \) is typically 5–10 MeV/c\(^2\) (1–2 MeV/c\(^2\)).

Electron candidates are required to satisfy tight PID criteria and to have lab-frame momentum \( |p_e| > 300 \) MeV/c, with an efficiency that rises from 85% at the lowest momenta to 95% for \( |p_e| > 1.0 \) GeV/c. Muon candidates are required to satisfy tight PID criteria; since muon PID relies on the hit pattern in the IFR, this effectively requires \( |p_\mu| \geq 600 \) MeV/c, and results in an efficiency of 40%–60% over the allowed momentum range. The energy of electron candidates is corrected for bremsstrahlung energy loss if photons are found close to the electron direction. Lepton candidates of either flavor are required to have at least 12 hits in the drift chamber and to have a laboratory-frame polar angle \( 0.4 < \theta < 2.6 \) rad (excluding the very forward and very backward regions of the tracking system) in order to ensure a well-measured momentum, since mis-measured lepton momenta distort the \( m^2_{miss} \) distribution and tend to move background events into the signal-like region. Approximately 5% of selected lepton candidates are misidentified, almost all of which are pions misconstructed as muons.

C. Total-Event and Single-Candidate Selection

We form whole-event candidates by combining \( B_{\text{tag}} \) candidates with \( D^{(*)}\ell^- \) candidate systems. We combine charged \( B_{\text{tag}} \) candidates with \( D^{(*)0}\ell^- \) systems and neutral \( B_{\text{tag}} \) candidates with both \( D^{(*)+}\ell^- \) and \( D^{(*)-}\ell^+ \) systems, where the inclusion of both charge combinations allows for neutral \( B \) mixing.

In correctly reconstructed signal and normalization events, all of the stable final-state particles, with the exception of the neutrinos, are associated with either the \( B_{\text{tag}} \), \( D^{(*)} \), or \( \ell^- \) candidate. Events with additional particles in the final state must therefore have been mis-reconstructed, and we suppress these backgrounds with two selection requirements on the “extra” particle content in the event. We require that all observed charged tracks be associated with either the \( B_{\text{tag}} \), \( D^{(*)} \), or \( \ell \) candidate. We compute \( E_{\text{extra}} \), the sum of the energies of all photon candidates not associated with the \( B_{\text{tag}} + D^{(*)}\ell \) candidate system, and we require \( E_{\text{extra}} < 150–300 \) MeV, depending on the \( D^{(*)} \) channel. When considering these extra tracks and extra photons, care is taken to reject track and photon candidates which are likely to be due to accelerator background, electronics noise, or reconstruction software failures; fake photons in the EMC are, to some degree, unavoidable, which is why we can not simply require \( E_{\text{extra}} = 0 \). The different \( D \) modes have very different levels of combinatorial background, which the \( E_{\text{extra}} \) cut is particularly effective at rejecting. Figure 5 shows distributions of \( E_{\text{extra}} \) for simulated signal and normalization events. Excellent agreement is seen in the two distributions, indicating that the efficiency of a cut on \( E_{\text{extra}} \) will largely cancel when we measure the branching-fraction ratio.

![FIG. 5: Distributions of \( E_{\text{extra}} \) in simulated events, shown after all event selection except the cut on \( E_{\text{extra}} \) itself. Signal \( B \to D^{(*)}\tau^- \gamma \) events are shown as points, while \( B \to D^{(*)}\ell^\pm \gamma \) normalization events are shown as the histogram. The gap between the \( E_{\text{extra}} = 0 \) bin and the remainder of the distribution corresponds to the minimum allowed photon energy, 50 MeV. The normalization is arbitrary. The agreement between the two distributions indicates that the efficiency of a cut on \( E_{\text{extra}} \) will cancel when we measure the branching-fraction ratio.](image-url)
4 (GeV/c\(^2\))^2, where \(q^2\) is calculated as

\[ q^2 = [p_B - p_{D(\pi)}]^2 = [p_{e^+e^-} - p_{tag} - p_{D(\pi)}]^2. \]  (10)

This requirement preferentially rejects combinatorial backgrounds from two-body B decays such as \(B \to D^{(*)}D\), where one D meson decays semileptonically (or, in the case of a \(D^+_s\), leptonically as \(D^+_s \to \pi^+ \nu\)). Our selection rejects about 25% of these backgrounds, while the signal efficiency is about 98% because signal events automatically satisfy \(q^2 > m^2_{\pi^+} \approx 3.16 \) (GeV/c\(^2\))^2. For \(B \to D\ell^-\tau^+\) decays, the \(q^2\) distribution peaks near zero (see Fig. 1), so this selection has an efficiency of about 60% for this normalization mode; for \(B \to D^{(*)}\ell^-\tau^+\) decays, the \(q^2\) distribution peaks at higher values, so our efficiency is about 70%. The \(q^2\) requirement is the main reason why the reconstruction efficiency is different for signal and normalization modes, as seen below.

If multiple candidate systems pass our selection in a given event, we select the one with the lowest value of \(E_{\text{extra}}\). This scheme preferentially selects the candidate that is least likely to have lost additional particles. The main effect of this algorithm is that a candidate in one of the \(D\ell^-\) channels will be selected before a candidate in one of the \(D\ell^+\) channels when both candidates are present in an event. Because \(D^* \to D\) feed-down is a dominant background while \(D \to D^*\) feed-up is comparatively rare, keeping as many true \(D^*\) events in the \(D^*\ell^-\) reconstruction channels helps to increase the sensitivity to the \(D\ell^-\) signals.

To improve the resolution on the missing momentum, we perform a kinematic fit to all \(\Upsilon(4S) \to \Lambda_{B(\pi)}D^{(*)}\ell^-\) candidates. We constrain charged track daughters of \(K^0_S\), D, and B mesons to originate from common vertices, and we constrain the \(\Upsilon(4S) \to B\overline{B}\) vertex to be consistent with the measured BABAR beamspot location. We constrain the mass of the signal D meson (and \(D^*\) meson, if there is one) to the measured value, and the combined momentum of the two B mesons to be consistent with the measured beam energy.

### D. \(D^{**}\) Control Sample Selection

We select four control samples to constrain the poorly known \(B \to D^{(*)}(\ell^-/\tau^-)\) background. The selection is identical to that of the signal channels, but we require the presence of a \(\pi^0\) meson in addition to the \(B_{\pi} + D^{(*)}\ell\) system. The \(\pi^0\) meson must have momentum greater than 400 MeV/c, and the event must satisfy \(E_{\text{extra}} < 500\) MeV, where the two photons from \(\pi^0 \to \gamma\gamma\) are excluded from the calculation of \(E_{\text{extra}}\).

Most of the \(D^{**}\) background in the four signal channels occurs when the \(\pi^0\) from \(D^{**} \to D^{(*)}\pi^0\) is not reconstructed, so these control samples provide a direct normalization of the background source. Similar \(D^{**}\) decays in which a \(\pi^\pm\) is lost contribute very little to the background since they do not have the correct charge correlation between the \(B_{\pi}\) and \(D^{(*)}\) candidate, and decays with two missing charged pions, which may have the correct charge correlation, have very low reconstruction efficiencies. The feed-down probabilities for the \(D^{**}(\ell^-/\tau^-)\) background are determined from simulation, with uncertainties in the \(D^{**}\) content treated as a systematic error as described in Sec. 1.13.

### V. SELECTED EVENT SAMPLES

After applying all of the criteria above, we select a total of 3196 data events, 2886 in the four signal channels and 310 in the \(D^{**}\) control samples, as listed in Table I. Since most of the events at large \(m^2_{\text{miss}}\) are either \(D\tau^-\) or \(D^+\tau^-\) signal events, the third column in the Table gives a first indication of where our sensitivity comes from. There are more events in the two \(B^-\) channels, \(D^0\ell^-\) and \(D^{*0}\ell^-\), due to a larger efficiency to reconstruct charged \(B_{\pi}\) candidates than neutral ones and, to a lesser extent, a larger efficiency to reconstruct \(D^0\) mesons on the signal side than \(D^+\) mesons. There are more events in the D channels than the \(D^*\) channels, particularly at large \(m^2_{\text{miss}}\), because these channels contain both D mesons and \(D^*\) feed-down. The greatest signal sensitivity therefore comes from the \(D^0\ell^-\) channel.

Figure 6 shows distributions of \(|p_\ell^0|\) versus \(m^2_{\text{miss}}\) for the selected data samples. One-dimensional distributions of \(m^2_{\text{miss}}\) and \(|p_\ell^0|\) for these samples are shown when we discuss the signal fit in Section VII.

Figure 7 shows distributions of \(|p_\ell|\) versus \(m^2_{\text{miss}}\) for several MC samples after applying all event selection criteria. While the composition of the event sample will be discussed in greater detail in the following section, these distributions exhibit the qualitative features of the data sample which are most relevant to our signal extraction. Figure 7(a) shows \(D^0\ell^-\) data, where we introduce the \(\Rightarrow\) notation to mean that these are true \(B^- \to D^0\ell^-\) events reconstructed in the \(D^0\ell^-\) channel. The \(m^2_{\text{miss}}\) distribution is very narrowly peaked around zero, as expected for one-neutrino events. Figure 7(b) shows \(D^{*0}\ell^-\) feed-down events where a \(D^{*0}\) is misreconstructed as a \(D^0\). In this case, the center of the \(m^2_{\text{miss}}\) distribution is offset from zero, and this offset decreases with increasing \(|p_\ell|\); this kinematic feature is common to all feed-down processes, and is due to

| Channel | \(N_{ev}\) | \(N_{ev} (large \ m^2_{\text{miss}})\) | \(N_{D^{**}\text{CS}}\) |
|---------|---------|-----------------|-----------------|
| \(D^0\ell^-\) | 1403 | 121 | 137 |
| \(D^{*0}\ell^-\) | 790 | 43 | 77 |
| \(D^+\ell^-\) | 295 | 36 | 66 |
| \(D^{*+}\ell^-\) | 398 | 14 | 30 |

**TABLE I:** Number of selected data events in the four signal channels, \(N_{ev}\), and in the four \(D^{**}\) control samples, \(N_{D^{**}\text{CS}}\). Here, the large \(m^2_{\text{miss}}\) region is taken to be \(m^2_{\text{miss}} > 2 \ (\text{GeV/c}^2)^2\) and corresponds to the region with greatest signal sensitivity.
to the fact that higher $|\mathbf{p}_\tau^*|$ correspond to lower $D^*$ momenta and therefore to lower momenta for the lost $\pi^0$ or $\gamma$. The width of the $m_{\text{miss}}^2$ distribution is also observed to decrease with increasing $|\mathbf{p}_\tau^*|$, a feature which is also common to most distributions; this narrowing is partly due to the same kinematic effect as before, the reduced $D^*$ phase space at large $|\mathbf{p}_\tau^*|$, and partly due to the fact that the lepton momentum resolution improves at higher momenta.

Figures 6(c) and (d) show similar distributions for signal events, (c) showing correctly reconstructed $B^- \rightarrow D^0 \tau^- \bar{\nu}_\tau$ and (d) showing $D^{*0}$ feed-down: in both plots, the large values of $m_{\text{miss}}^2$ due to the three neutrinos are clearly observed. Again, the $m_{\text{miss}}^2$ distributions move towards zero and become narrower at high $|\mathbf{p}_\tau^*|$, in this case due to the reduced phase space for the multiple neutrinos, although, in Fig. 6(d), the effect of the lost $\pi^0$ or $\gamma$ can also be seen as a deficit along the lower-left edge of the distribution. Figure 6(e) shows feed-down from $B \rightarrow D^{**} \ell^- \bar{\nu}_\ell$ into the $D^0 \ell^-$ channel, where, in addition to the neutrino, one or more $\pi^0$ mesons or photons from the $D^{**}$ decay have been lost. Since $\pi^0$ mesons from $D^{**}$ decay typically have higher momentum than those from $D^*$ decay, the $m_{\text{miss}}^2$ distribution is much broader than that in Fig. 6(b). Figure 6(f) shows the feed-up process $D^0 \ell^- \bar{\nu}_\ell \Rightarrow D^0 \ell^- \gamma_\nu$, where a true $D^0$ meson is paired with a combinatorial $\pi^0$ or $\gamma$ to fake a $D^{*0}$ candidate. In this case, the $m_{\text{miss}}^2$ distribution is shifted in the opposite direction from Fig. 6(b). Figures 6(g), (h), and (i) show three additional distributions for events re-constructed in the $D^{*0} \ell^-$ channel, $D^{*0} \ell^- \bar{\nu}_\ell$, $D^{*0} \tau^- \bar{\nu}_\tau$ signal, and $D^{**} \ell^- \bar{\nu}_\ell$ background, respectively; each of these distributions is similar to the corresponding one in the $D^0 \ell^-$ channel, Figs. 6(a), (c), and (e), respectively.

Figures 6(j) and (k) show charge-crossfeed backgrounds: true $B \rightarrow D^{(*)} \ell^- \bar{\nu}_\ell$ events reconstructed with the wrong charge for both the $B_{\text{tag}}$ and $D^{(*)}$ meson. Typically this occurs when a low-momentum $\pi^\pm$ is swapped between the two mesons. Note that, even though the event is misreconstructed, this particle misassignment does not substantially alter the total missing momentum, so that the $m_{\text{miss}}^2$ distribution still peaks at or near zero. While the events in Fig. 6(k), which are reconstructed in the $D^{*} \ell^-$ channels, are very strongly peaked at $m_{\text{miss}}^2 = 0$, Fig. 6(j) includes a large feed-down component, and therefore exhibits the same sloping behavior seen in Fig. 6(b).

Figure 6(l) shows the distribution for combinatorial background for all four signal channels. This background is dominated by hadronic $B$ decays such as $B \rightarrow D^{(*)} D^{(*)}_{\text{tag}}$ that produce a secondary lepton, including events with $\tau$ leptons from $D_s$ decay.

In our $B\bar{B}$ MC sample, our criteria select $D^{**}$ control samples which are 60%–80% pure $B \rightarrow D^{**} \ell^- \bar{\nu}_\ell$ events, of which more than 90% involve true $D^{**} \rightarrow D^{(*)} \pi^0$ transitions. The remaining events are split between feed-up from $B \rightarrow D^{(*)} \ell^- \bar{\nu}_\ell$ and combinatorial background. In these control samples, the $B \rightarrow D^{**} \ell^- \bar{\nu}_\ell$ component peaks at or near zero in $m_{\text{miss}}^2$, just as $B \rightarrow D^{(*)} \ell^- \bar{\nu}_\ell$ does in the four signal channels. The qualitative features of the other contributions are similar to what is seen in the signal channels.

VI. KINEMATIC CONTROL SAMPLES

The event selection criteria described in Section IV are more complicated than those used in a typical $B\bar{B}$ analysis, due to the full-event reconstruction of a high-multiplicity final state and the need to veto events with extra tracks and neutral clusters. We use two data control samples to validate our simulation with respect to the observed behavior in data. The control samples are
kinematically selected, with no requirement on \( m_{\text{miss}}^2 \), to be high purity samples of \( B \to D^{(*)}\ell\bar{\nu}_\ell \) events, with little or no contamination from signal decays.

The first control sample is defined by requiring the reconstructed lepton to satisfy \( |p_\ell^*| > 1.5 \text{ GeV}/c \), and is therefore a subset of our full analysis sample. In simulation, 95% of the selected sample is \( B \to D^{(*)}\ell\bar{\nu}_\ell \) (the two \( D^*\ell^- \) channels are approximately 95% \( B \to D^*\ell^-\bar{\nu}_\ell \), while the two \( D\ell^- \) channels include both \( B \to D\ell^-\bar{\nu}_\ell \) and large feed-down from \( B \to D^{(*)}\ell\bar{\nu}_\ell \). The remaining 5% of the sample is composed of about 1%-3% \( B \to D^{(*)}\ell^-\bar{\nu}_\ell \), less than 1% \( B \to D^{(*)}\tau^-\bar{\nu}_\tau \), and 1%-2% of combinatorial backgrounds.

For the second control sample, we remove the standard \( q^2 > 4 \text{ (GeV}/c^2)^2 \) selection and require that events instead satisfy \( q^2 < 5 \text{ (GeV}/c^2)^2 \), with \( q^2 \) calculated according to Eq. (10). This control sample has very little overlap with our final event sample, where we require \( q^2 > 4 \text{ (GeV}/c^2)^2 \). Although the two control samples do have some overlap, this \( q^2 \) control sample has the advantage over the first of allowing us to examine events with low \( |p_\ell^*| \), as expected for signal events. In simulation, approximately 90% of this control sample is \( B \to D^{(*)}\ell\bar{\nu}_\ell \) (as in the first sample, the two \( D^*\ell^- \) channels are approximately 90% \( B \to D^*\ell^-\bar{\nu}_\ell \), while the two \( D\ell^- \) channels include \( D^* \) feed-down). The remainder of the sample is composed of about 3% \( B \to D^{(*)}\ell^-\bar{\nu}_\ell \), 3% \( B \to D^{(*)}\tau^-\bar{\nu}_\tau \), and 4%-5% combinatorial backgrounds.

Figure 7 shows several data-simulation comparisons in the two control samples. The four \( D^{(*)}\ell^- \) channels have been combined in these plots as have the two control samples, and this union of the two control samples is responsible for the large steps visible in (a) and (b). We see good agreement between data and simulation in these plots, as well as in similar studies where the two control samples are examined separately, the four \( D^{(*)}\ell^- \) channels are examined separately, the two lepton types are examined separately, and where the data are split according to \( B\bar{B} \) running period. We have examined variables related to \( B_{\text{tag}} \) reconstruction, signal-side reconstruction, hermeticity and whole-event reconstruction, and missing momentum. In all cases, we observe that the simulation does a reasonable job describing the data. Because of the relative normalization scheme, small differences between simulation and data have no detrimental effect on the analysis.

VII. FIT OF SIGNAL AND BACKGROUND YIELDS

A. Fit Overview

Signal and background yields are extracted using an extended unbinned maximum likelihood fit to the joint \((m_{\text{miss}}^2, |p_\ell^*|)\) distribution. The fit is performed simultaneously in the four signal channels and the four \( D^* \) control samples. Two two-dimensional probability density

![Figure 7](https://via.placeholder.com/150)
functions (PDFs) are presented in Section VII B; each component in the fit (listed below) is described by one of these two PDFs, with parameters determined from fits to simulated event samples. A set of constraints, described in Section VII B, relate fit components in different reconstruction channels. These constraints are also determined from MC samples, except for parameters describing the amount of $D^*$ feed-down into the $D\ell^-$ signal channels, which are determined directly by the fit to data.

Tables II and III summarize the parameterization of the fit in the four signal channels and the four $D^{*+}$ control samples, respectively. In each of the four signal channels, we describe the data as the sum of seven components: $D\tau, D^{*+}\tau\tau, D\ell^-\tau\pi, D^{*+}\ell^-\tau\pi, D^{(*)}\ell^-\tau\pi, D^{(*)}\ell^-\tau\pi, D^{(*)}\ell^-\pi\pi$, charge-crossfeed, and combinatorial background. Each of these components is described by one of the two PDFs given in Section VII B with the numerical parameters of the 32 PDFs determined from independent MC samples. The charge-crossfeed components in the two $D\ell^-$ signal channels are described by a single PDF, with common parameters for $D^0\ell^-$ and $D^+\ell^-$, as are the two $D^{*}\ell^-$ charge-crossfeed components and the four $D^{(*)}\ell^-\pi^0$ components; the four combinatorial background components in the signal channels are described by a single PDF with common parameters, as are the four in the $D^{*+}$ control samples.

$B\rightarrow D^{(*)}\ell^-\tau\pi$, events feeding up into the $D^{*+}$ control samples are expected to contribute $1.8 \pm 0.6$ events in the four channels together, so these events are combined with the light lepton contribution. In both the control samples and in the signal channels, $B\rightarrow D^{*+}\tau\pi$, events...
and charge-crossfeed backgrounds in both the signal channels. The fit to the expected value from simulation, as are the background yields in the four signal channels are fixed in modes and one for neutral is treated as a systematic uncertainty below.

| Channel Source | Abundance in BFB MC (%) |
|----------------|------------------------|
| D^*0 \ell^-   | D^0 \tau^- \pi^+ signal     | 5   |
|                | D^0 \tau^- \pi^+ signal feed-up | 0.5 |
|                | D^*0 \ell^- \tau^+ normalization | 90  |
|                | D^0 \ell^- \tau^+ feed-up     | 2   |
|                | D^*0 (\ell^-/\tau^-) \pi^- feed-down | 3   |
|                | Charge-crossfeed             | 0.5 |
|                | Combinatorial background    | 1   |
|                | D^*0 \tau^- \ell^- signal   | 3   |
|                | D^0 \tau^- \ell^- signal feed-up | 3   |
|                | D^0 \ell^- \tau^- normalization | 25  |
|                | D^0 \ell^- \tau^- feed-down | 60  |
|                | D^*0 (\ell^-/\tau^-) \pi^- feed-down | 2   |
|                | Charge-crossfeed             | 2   |
|                | Combinatorial background    | 2   |
|                | D^*+ \tau^- \ell^- signal   | 5   |
|                | D^+ \tau^- \ell^- signal feed-up | 0.1 |
|                | D^+ \ell^- \tau^- normalization | 90  |
|                | D^+ \ell^- \tau^- feed-up   | 0.5 |
|                | D^*+ (\ell^-/\tau^-) \pi^- feed-down | 3   |
|                | Charge-crossfeed             | 0.1 |
|                | Combinatorial background    | 2   |
|                | D^+ \tau^- \ell^- signal    | 5   |
|                | D^+ \tau^- \ell^- signal feed-down | 2   |
|                | D^+ \ell^- \tau^- normalization | 45  |
|                | D^+ \ell^- \tau^- feed-down | 40  |
|                | D^*+ (\ell^-/\tau^-) \pi^- feed-down | 6   |
|                | Charge-crossfeed             | 1   |
|                | Combinatorial background    | 3   |

TABLE III: Components of the signal extraction fit in the D^{**} control sample channels, and their approximate abundances in our BFB MC sample. The structure of the fit is identical between the B^- and B^0 channels. There are five components in each of the four D^{**} control sample channels.

| Channel Source | Abundance in BFB MC (%) |
|----------------|------------------------|
| D^{**0} \ell^- | D^{**+} \tau^- \pi^+ signal | 60  |
|                | D^{**+} \tau^- \pi^+ signal | 18  |
|                | D^{**0} \ell^- \tau^- feed-up | 2   |
|                | D^{**+} \ell^- \tau^- feed-up | 3   |
|                | Charge-crossfeed          | 1   |
|                | Combinatorial background  | 20  |
| D^{**0} \tau^- \ell^- | D^{**+} \tau^- \pi^+ signal | 70  |
|                | D^{**+} \tau^- \pi^- feed-up | 10  |
|                | D^{**0} \ell^- \tau^- feed-up | 3   |
|                | Charge-crossfeed          | 2   |
|                | Combinatorial background  | 15  |
| D^{**+} \tau^- \ell^- | D^{**+} \tau^- \pi^+ signal | 65  |
|                | D^{**+} \ell^- \tau^- feed-up | 20  |
|                | D^{**+} \ell^- \tau^- feed-up | 0.1 |
|                | Charge-crossfeed          | 0.1 |
|                | Combinatorial background  | 15  |
| D^{**+} \tau^- \ell^- | D^{**+} \tau^- \pi^+ signal | 60  |
|                | D^{**+} \ell^- \tau^- feed-up | 6   |
|                | D^{**+} \ell^- \tau^- feed-up | 3   |
|                | Charge-crossfeed          | 1   |
|                | Combinatorial background  | 30  |

are expected to contribute 3.5%-4.5% of the total D^{**} yield; these events are combined with the light lepton contribution, and the amount of D^{**-}\tau^- is varied as a systematic uncertainty.

The fit has 18 free parameters: four signal branching-fraction ratios R, one for each D^{(*)} meson; four B \rightarrow D^{(*)} \ell^- \tau^- normalization yields; four B \rightarrow D^{**} \ell^- \tau^- background yields; four combinatorial background yields, one in each of the four D^{**} control samples; two parameters describing D^* \Rightarrow D feed-down, one for charged B modes and one for neutral B modes. The combinatorial background yields in the four signal channels are fixed in the fit to the expected value from simulation, as are the charge-crossfeed backgrounds in both the signal channels and D^{**} control samples; variation of these backgrounds is treated as a systematic uncertainty below.

We also perform a second, B^-\bar{B}^0 constrained, fit, by requiring R(D^+) = R(D^0) and R(D^{**}) = R(D^{**0}), reducing the number of free parameters to 16.

B. Probability Density Functions

We construct an empirical model of the two-dimensional (m_{\text{miss}}^2, |p_T^\ell|) PDF as the product of two terms: a one-dimensional function to describe the |p_T^\ell| distribution, discussed in Section VII.B.1 and a |p_T^\ell|-dependent “resolution” function to describe the m_{\text{miss}} distribution, to be discussed in Section VII.B.2. For processes in which the only missing particle is a single neutrino, the true m_{\text{miss}} spectrum is a delta function located at zero and the observed distribution is a pure resolution function. For components with multiple missing particles, the observed m_{\text{miss}} distribution is the convolution of the physical m_{\text{miss}} spectrum with our detector resolution. The PDFs presented below are used to describe both of these physical cases, with different numerical parameters describing the different behaviors; these two PDFs are flexible enough to describe the variety of physical and

7 This constraint follows from isospin symmetry in both the signal and normalization modes but is more general.
resolution processes needed in this analysis.

1. One-dimensional \(|p_\ell^*|\) Parameterization

We use a generalized form of a Gaussian to model the \(|p_\ell^*|\) distribution. The Gaussian distribution,

\[
\mathcal{G}(|p_\ell^*|) \propto \exp \left( -\frac{1}{2} \left| \frac{|p_\ell^*| - p_0}{\sigma} \right|^2 \right),
\]

(11)

has the same general properties as our distributions: it rises smoothly from zero to a peak value and then falls smoothly back to zero again. Here, \(p_0\) represents the value of \(|p_\ell^*|\) for which \(\mathcal{G}\) peaks and \(\sigma\) represents the width of the Gaussian distribution.

This gross agreement is not enough, however, so we define a modified Gaussian function,

\[
\mathcal{H}(|p_\ell^*|) \propto \exp \left( -\frac{1}{2} \left| \frac{|p_\ell^*| - p_0}{\sigma(|p_\ell^*|)} \right|^\nu(|p_\ell^*|) \right),
\]

(12)

where, for convenience, we have absorbed the constant factor of 2 into the definition of \(\sigma(|p_\ell^*|)\). By allowing the width and exponent of the Gaussian to be functions of \(|p_\ell^*|\), we are able to describe a greater variety of shapes. Specifically, we take \(\nu(|p_\ell^*|)\) to be a linear function,

\[
\nu(|p_\ell^*|) = \nu_L + \frac{\nu_H - \nu_L}{2.4 \text{ GeV}/c} \cdot |p_\ell^*|
\]

(13)

where \(\nu_L\) and \(\nu_H\) are the values of the exponential term at the low and high endpoints of \(|p_\ell^*|\), fixed at zero and 2.4 GeV/c, respectively. Similarly, we parameterize \(\sigma(|p_\ell^*|)\) as a bilinear function,

\[
\sigma(|p_\ell^*|) = \begin{cases} 
\sigma_L + \frac{\sigma_0 - \sigma_L}{p_0} \cdot |p_\ell^*| & |p_\ell^*| < p_0 \\
\sigma_0 + \frac{\sigma_L - \sigma_0}{2.4 \text{ GeV}/c - p_0} \cdot (|p_\ell^*| - p_0) & |p_\ell^*| > p_0
\end{cases}
\]

(14)

where \(\sigma_L\), \(\sigma_0\), and \(\sigma_H\) represent the widths of the Gaussian at \(|p_\ell^*| = 0\), \(|p_\ell^*| = p_0\), and \(|p_\ell^*| = 2.4 \text{ GeV}/c\), respectively. Even though this parameterization is discontinuous at the point \(|p_\ell^*| = p_0\), the resulting function \(\mathcal{H}(|p_\ell^*|)\) remains smooth since the numerator in the exponent, \(|p_\ell^*| - p_0\), goes to zero at the same point.

The \(|p_\ell^*|\) parameterization therefore has six free parameters: \(p_0\), the peak; \(\nu_L\) and \(\nu_H\), describing the exponential term; and \(\sigma_L\), \(\sigma_0\), and \(\sigma_H\), describing the width. When performing fits using this PDF, we integrate \(\mathcal{H}\) numerically to compute the normalization.

2. Two-dimensional PDF Parameterization

We construct two types of two-dimensional PDF, \(\mathcal{P}_1(|p_\ell^*, m_{\text{miss}}^2|)\) and \(\mathcal{P}_2(|p_\ell^*, m_{\text{miss}}^2|)\) by multiplying the model of the lepton spectrum above by a “resolution” function in \(m_{\text{miss}}^2\), where the resolution is a function of \(|p_\ell^*|\). Allowing the parameters of the resolution function to be functions of \(|p_\ell^*|\) produces a correlation between the two fit variables, and it is these parameters which allow the PDFs to describe such a wide variety of shapes.

Using the model of the lepton spectrum \(\mathcal{H}(|p_\ell^*|)\) introduced above, we construct the PDFs as:

\[
\mathcal{P}_1(|p_\ell^*, m_{\text{miss}}^2|) \equiv \mathcal{H}(|p_\ell^*|) \times \left[ f_1(|p_\ell^*|) \mathcal{G}_1(|p_\ell^*|; m_{\text{miss}}^2) + (1 - f_1(|p_\ell^*|)) \mathcal{G}_2(|p_\ell^*|; m_{\text{miss}}^2) \right]
\]

(15)

and

\[
\mathcal{P}_2(|p_\ell^*, m_{\text{miss}}^2|) \equiv \mathcal{H}(|p_\ell^*|) \times \left[ f_1(|p_\ell^*|) \mathcal{G}_1(|p_\ell^*|; m_{\text{miss}}^2) + f_2(|p_\ell^*|) \mathcal{G}_b(|p_\ell^*|; m_{\text{miss}}^2) + (1 - f_1(|p_\ell^*|) - f_2(|p_\ell^*|)) \mathcal{G}_2(|p_\ell^*|; m_{\text{miss}}^2) \right].
\]

(16)

Here, the functions \(\mathcal{G}_1\) and \(\mathcal{G}_2\) are Gaussians and \(\mathcal{G}_b\) is a bifurcated Gaussian (Gaussian with different \(\sigma\) parameters on either side of the mean), respectively; all are functions of \(m_{\text{miss}}^2\), with parameters dependent on \(|p_\ell^*|\).

The \(|p_\ell^*|\) dependence of the various parameters of \(\mathcal{G}_1, \mathcal{G}_2\) and \(\mathcal{G}_b\) is listed in Table [X]. The total number of free parameters for \(\mathcal{P}_1\) is 18: six for \(\mathcal{H}(|p_\ell^*|)\), five each for \(\mathcal{G}_1\) and \(\mathcal{G}_2\), and two each for \(f_1\) and \(f_2\). The total number of free parameters for \(\mathcal{P}_2\) is 24: six for \(\mathcal{H}(|p_\ell^*|)\), five each for \(\mathcal{G}_1\) and \(\mathcal{G}_2\), four for \(\mathcal{G}_b\), and two each for \(f_1\) and \(f_2\).

We use the simpler PDF, \(\mathcal{P}_1\), to model most of the semileptonic fit components (22 out of 32), as well as the charge-crossfeed and combinatorial backgrounds. For the remaining ten components, however, the more complicated parameterization \(\mathcal{P}_2\) is required to adequately describe the \(m_{\text{miss}}^2\) tail. Eight of these components are the ones in which the only missing particle is a single neutrino,

\[
D^{\mu\nu} - \mathcal{T}_\ell \Rightarrow D^{\mu\nu} \ell^- \quad D^{\nu\mu} - \mathcal{T}_\ell \Rightarrow D^{\nu\mu} \ell^- \\
D^{\mu\nu} - \mathcal{T}_\ell \Rightarrow D^{\mu\nu} \pi^0 \ell^- \quad D^{\nu\mu} - \mathcal{T}_\ell \Rightarrow D^{\nu\mu} \pi^0 \ell^- \\
D^{\mu\nu} - \mathcal{T}_\ell \Rightarrow D^{\mu\nu} \gamma \ell^- \quad D^{\nu\mu} - \mathcal{T}_\ell \Rightarrow D^{\nu\mu} \gamma \ell^- \\
D^{\mu\nu} - \mathcal{T}_\ell \Rightarrow D^{\mu\nu} \tau \ell^- \quad D^{\nu\mu} - \mathcal{T}_\ell \Rightarrow D^{\nu\mu} \tau \ell^- \\
\]

and the remaining two are components in which a single neutrino and a soft \(\pi^0\) or \(\gamma\) are missing,

\[
D^{\mu\nu} - \mathcal{T}_\ell \Rightarrow D^{\mu\nu} \ell^- \quad D^{\nu\mu} - \mathcal{T}_\ell \Rightarrow D^{\nu\mu} \ell^- .
\]
TABLE IV: $|p_t^*|$ dependence of the $m^2_{\text{miss}}$ PDF parameterization. The form of $f_2$ is chosen to allow the $G_b$ term to contribute at low $|p_t^*|$, but to drive this term rapidly to zero as $|p_t^*|$ increases. The form of $\sigma_H$ is chosen to allow for a long tail towards high $m^2_{\text{miss}}$ at low $|p_t^*|$, but to drive this term rapidly to zero as $|p_t^*|$ increases (note that there is no problem having $\sigma$ approach zero since the amplitude of this term goes to zero as well; the result is finite and well-behaved). $N_{\text{par}}$ gives the number of free parameters for each term separately.

| Function Parameter Dependence on $|p_t^*|$ | $N_{\text{par}}$ |
|------------------------------------------|----------------|
| $G_{1,2}$ mean quadratic | 3 |
| $G_{1,2}$ $\sigma$ linear | 2 |
| $G_b$ mean constant | 1 |
| $G_b$ $\sigma_L$ constant | 1 |
| $G_b$ $\sigma_H$ $\sigma_{H0} \cdot \left[1 - \left(\frac{|p_t^*|}{2.4 \, \text{GeV}/c}\right)^2\right]$ | 2 |
| $P_{1,2}$ $f_1$ linear | 2 |
| $P_2$ $f_2$ $F \left(\frac{2.4 \, \text{GeV}/c - |p_t^*|}{2.2 \, \text{GeV}/c}\right)^{\delta}$ | 2 |

C. Crossfeed Constraints

We apply a number of constraints in the fit, relating the event yields between different reconstruction channels in order to make use of all available information. These constraints help to maximize our sensitivity, particularly to the $B \to D\tau^-\bar{\nu}_\tau$ signals where the dominant backgrounds are due to feed-down. There are 20 such constraints in the fit, corresponding to 20 different ways in which a true $B \to D/D^*/D'^* \ell^- \bar{\nu}_\tau$ event can be reconstructed with the wrong final-state meson, either as feed-down ($D^* \to D$ and $D'^* \to D/D^*$) or as feed-up ($D \to D^*\pi$ and $D^* \to D^*\pi^0$).

These constraints are implemented in the fit by requiring that the number of events of type $j$ correctly reconstructed in the $i$th channel ($N_{ij}$) is related to the number of events of type $j$ reconstructed in a crossfeed channel $i'$ ($N_{i'j}$) by

$$N_{i'j} = N_{ij} \cdot f_{i \to i'j} \times$$

where $f_{i \to i'j}$ is a crossfeed constraint relating the two yields. The crossfeed constraints $f_{i \to i'j}$ are linearly related to the misreconstruction probability. For feed-down processes, in which the probability to lose a low-momentum $\pi^0$ or $\gamma$ is high, $f_{i \to i'j}$ typically takes values between 0.2 and 1.0; for feed-up processes, in which the probability to reconstruct a fake $\pi^0$ or $\gamma$ in a narrow mass window is low, $f_{i \to i'j}$ typically takes values between 0.01 and 0.1.

The values for most of the $f_{i \to i'j}$ terms are taken from simulation, but, in order to reduce systematic effects, the values of the dominant feed-down components, $B \to D^* \ell^- \bar{\nu}_\tau$ reconstructed in the $D\ell^-$ signal channels, are left free in the fit to data. We also use the floating values of these $D^*$ feed-down constraints to apply a small first-order correction to the corresponding signal feed-down constraints describing $B \to D^+ \tau^- \bar{\nu}_\tau$ reconstructed in the $D\ell^-$ channels; in this way, we use the high-statistics $D^+ \ell^- \bar{\nu}_\tau$ samples to improve our knowledge of the signal feed-down probability.

D. Projections of the Fit to Data

Figures 10–14 show projections of the $B^-\bar{D}^0$-constrained fit. Figure 10 shows projections in $m^2_{\text{miss}}$ for the four signal channels, showing both the low $m^2_{\text{miss}}$ region, which is dominated by the normalization modes $B \to D^\pm \ell^- \bar{\nu}_\tau$, and the high $m^2_{\text{miss}}$ region, which is dom-
FIG. 10: (Color online) Distributions of events and fit projections in $m_{\text{miss}}^2$ for the four final states: $D^{*0}\ell^-$, $D^0\ell^-$, $D^{*+}\ell^-$, and $D^+\ell^-$. The normalization region $m_{\text{miss}}^2 \approx 0$ is shown with finer binning in the insets. The fit components are combinatorial background (white, below dashed line), charge-crossfeed background (white, above dashed line), the $B \rightarrow D\ell^-$ normalization mode (/// hatching, yellow), the $B \rightarrow D^*\ell^-$ normalization mode (\ hatching, light blue), $B \rightarrow D^{*}\ell^-$ background (dark, or blue), the $B \rightarrow D\tau^-\nu_\tau$ signal (light grey, green), and the $B \rightarrow D^{*}\tau^-\nu_\tau$ signal (medium grey, magenta). The fit shown incorporates the $B^- - B^0$ constraints.

FIG. 11: (Color online) Distributions of events and fit projections in $|p_\tau^*|$ for the four final states $D^{*0}\ell^-$, $D^0\ell^-$, $D^{*+}\ell^-$, and $D^+\ell^-$, shown in the normalization region, $m_{\text{miss}}^2 < 1 (\text{GeV}/c^2)^2$. The fit components are shaded as in Fig. 10.

VIII. SIGNAL EXTRACTION AND NORMALIZATION

The fit described in Section VII directly measures, for each signal mode, the ratio of the number of signal events in the data sample, $N_{\text{sig}}$, to the number of corresponding normalization events, $N_{\text{norm}}$. We measure the signal branching-fraction ratios $R$ as

$$R \equiv \frac{N_{\text{sig}}}{N_{\text{norm}}} \cdot \frac{1}{\varepsilon_{\text{sig}}/\varepsilon_{\text{norm}}} \cdot \frac{1}{B(\tau^- \rightarrow \ell^- \nu_\tau \nu_\tau)},$$

where the relative efficiency $\varepsilon_{\text{sig}}/\varepsilon_{\text{norm}}$ is calculated from signal MC samples as...
FIG. 12: (Color online) Distributions of events and fit projections in $|p\star\ell|$ for the four final states $D^\star 0 \ell^-$, $D^0 \ell^-$, $D^{+\star} \ell^-$, and $D^{+} \ell^-$, shown in the signal region, $m^{miss}_2 > 1$ (GeV/c$^2$)$^2$. The fit components are shaded as in Fig. 10.

Signal efficiencies are given in Table VI. The relative efficiencies for the two $B \to D\tau^-\bar\nu_\tau$ modes are much larger than unity because of the $q^2$ cut, which is $\approx 98\%$ efficient for signal events but rejects about 50% of the $B \to D\ell^-\bar\nu_\ell$ normalization events, as seen in Fig. 1(a). The $q^2$ cut has a similar, but less pronounced, effect on the $D^*$ modes, but, due to the lower efficiency for identifying secondary leptons in the signal modes, the resulting relative efficiency is close to unity.

IX. SYSTEMATIC UNCERTAINTIES

Table VII summarizes all of the systematic uncertainties considered in this analysis. Because our signal is ex-
tracted and normalized relative to $B \to D^{(*)}\ell^{-}\tau_{\ell}^-$, many sources of systematic uncertainty—especially those related to reconstruction efficiency—are expected to cancel, either partially or completely, when we take the ratio.

We describe the individual contributions to the systematic uncertainty below. We divide the systematics into two broad categories: additive and multiplicative. Additive systematic uncertainties are those which affect the fit yields and therefore reduce the significance of the measured signals. Multiplicative uncertainties affect the normalization of the signals and the numerical results but not the significance.

A. Additive Systematic Uncertainties

In order to estimate additive systematic uncertainties, we perform an ensemble of fits to MC event samples. For each source of uncertainty, we perform a number of tests where we modify, as appropriate, the fit shapes, crossfeed constraints, and the combinatorial background yields (all of which are fixed to MC-derived values in the nominal fit) and perform a signal fit. By doing a large number of such tests and studying the distribution of fit results in these ensembles, we are able to estimate the systematic uncertainties. In all of these ensembles, we take the RMS of the observed distribution, relative to the corresponding mean fit value, as the systematic uncertainty.

1. Monte Carlo Statistics

In order to study the systematic uncertainties due to limited Monte Carlo statistics, we perform two ensembles of fits. In the first ensemble, we perform a variation of the PDF shapes. Each of the 37 PDFs are independently varied by generating new values for each of the 18 or 24 shape parameters according to the uncertainties in the PDF fit, taking into account correlations between the fitted parameters. In the second ensemble, we vary each of the feed-up and feed-down constraints according to their statistical uncertainties.

Figure 15 shows distributions of fit results for the ensemble of PDF shape fits.

2. Combinatorial Background Modeling

Table VII summarizes the physical sources of combinatorial background considered in this analysis, including their approximate abundances in our $B\bar{B}$ MC sample after all event selection. In order to study systematic effects, we perform an ensemble of fits, reweighting events from the various combinatorial sources.

In total, the two-body $B$ decays $B \to D^{(*)}D^{(*)}$ and $B \to D^{(*)}D^{(*)}$ constitute approximately 45% of the total combinatorial background yield, while the three-body decays $B \to D^{(*)}D^{(*)}K$ contribute another 15%.

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**FIG. 14:** (Color online) Distributions of events and fit projections in $|p_\ell^*|$ for the four $D^{(*)}$ control samples $D^0\pi^0\ell^-$, $D^0\pi^0\ell^-$, $D^+\pi^0\ell^-$, and $D^+\pi^0\ell^-$. The fit components are shaded as in Fig. 10.

**TABLE V:** Relative signal efficiencies $\varepsilon_{\text{sig}}/\varepsilon_{\text{norm}}$ for the four signal modes.

| Signal mode       | $\varepsilon_{\text{sig}}/\varepsilon_{\text{norm}}$ |
|-------------------|-----------------------------------------------|
| $B^- \to D^0\tau^-\pi_\ell^-$ | $1.85 \pm 0.02$ |
| $B^- \to D^{*0}\tau^-\pi_\ell^-$ | $0.99 \pm 0.01$ |
| $B^0 \to D^+\tau^-\pi_\ell^-$ | $1.83 \pm 0.03$ |
| $B^0 \to D^{*+}\tau^-\pi_\ell^-$ | $0.91 \pm 0.01$ |
TABLE VI: Contributions to the total systematic uncertainty. The additive systematic uncertainties represent uncertainties on the fit yield, and therefore reduce the statistical significance of the results. The multiplicative systematic uncertainties represent uncertainties on the normalization, so they affect the numerical results but not the statistical significance. The first four columns summarize errors on the individual branching-fraction ratios; the last two columns summarize errors on the $B^–B^0$ constrained measurement. The totals here refer to errors on the branching-fraction ratios $R$; the errors on $B(B \to D^{(*)}\ell^+\nu_\ell)$ (discussed in Section X) only apply to the absolute branching fractions, and are not included in the quoted total error.

| Source                                      | $D^0\tau\nu$ | $D^{*0}\tau\nu$ | $D^+\tau\nu$ | $D^{*+}\tau\nu$ | $D^+\pi\nu$ | $D^{*+}\pi\nu$ |
|---------------------------------------------|---------------|------------------|---------------|------------------|--------------|---------------|
| MC stat. (PDF shape)                        | 11.5          | 8.4              | 4.5           | 1.8              | 6.9          | 4.7           |
| MC stat. (constraints)                      | 4.2           | 1.9              | 6.1           | 1.3              | 3.6          | 1.4           |
| Comb. BG modeling                            | 7.5           | 4.1              | 11.5          | 2.6              | 9.1          | 2.9           |
| $D^{*+}$ modeling                           | 5.7           | 0.5              | 1.6           | 0.2              | 3.0          | 0.4           |
| $B\to D^*$ form factors                     | 1.9           | 0.7              | 0.8           | 0.2              | 1.4          | 0.4           |
| $B\to D$ form factors                       | 0.2           | 0.7              | 0.6           | 0.2              | 0.3          | 0.4           |
| $m^2_{\text{miss}}$ tail modeling           | 1.5           | 0.5              | 1.2           | 0.4              | 1.6          | 0.1           |
| $\pi^0$ crossfeed constraints               | 0.5           | 1.1              | 0.5           | 0.9              | 0.5          | 1.0           |
| $D^{*+}$ feed-down                          | 0.4           | 0.1              | 0.1           | 0.3              | 0.2          | 0.2           |
| $D^{*+}\tau^-\nu_\tau$ abundance           | 0.4           | 1.3              | 0.3           | 0.2              | 0.3          | 0.8           |
| Total additive                               | 15.6          | 9.7              | 14.0          | 3.6              | 12.5         | 5.8           |

Additive systematic uncertainties

| Source                                      | $D^0\tau\nu$ | $D^{*0}\tau\nu$ | $D^+\tau\nu$ | $D^{*+}\tau\nu$ | $D^+\pi\nu$ | $D^{*+}\pi\nu$ |
|---------------------------------------------|---------------|------------------|---------------|------------------|--------------|---------------|
| MC stat. (efficiency)                       | 1.2           | 1.1              | 1.5           | 1.1              | 1.0          | 0.8           |
| Bremsstrahlung/FSR                          | 0.6           | 0.5              | 0.3           | 0.4              | 0.4          | 0.5           |
| Tracking $\varepsilon$                      | 0.0           | 0.0              | 0.0           | 0.0              | 0.0          | 0.0           |
| $e$ PID $\varepsilon$                       | 0.5           | 0.5              | 0.6           | 0.6              | 0.6          | 0.6           |
| $\mu$ PID $\varepsilon$                     | 0.5           | 0.6              | 0.6           | 0.6              | 0.6          | 0.6           |
| $K$ PID $\varepsilon$                       | 0.2           | 0.1              | 0.2           | 0.0              | 0.2          | 0.0           |
| $\pi$ PID $\varepsilon$                     | 0.1           | 0.1              | 0.2           | 0.0              | 0.1          | 0.1           |
| $K^0_{\Sigma}$ $\varepsilon$                | 0.1           | 0.0              | 0.1           | 0.1              | 0.1          | 0.1           |
| Neutral ($\pi^0$ and $\gamma$) $\varepsilon$| 0.0           | 0.0              | 0.0           | 0.1              | 0.0          | 0.0           |
| Daughter $B$’s                               | 0.1           | 0.3              | 0.0           | 0.1              | 0.1          | 0.3           |
| $B(\tau^-\to \ell^-\nu_\ell)$              | 0.2           | 0.2              | 0.2           | 0.2              | 0.2          | 0.2           |
| Total multiplicative                         | 15.6          | 9.9              | 14.0          | 3.9              | 12.5         | 6.0           |
| $B(B \to D^{(*)}\ell^-\nu)$                | 10.2          | 7.7              | 9.4           | 3.7              | 6.8          | 3.4           |

Multiplicative systematic uncertainties

TABLE VII: Sources and approximate abundances of combinatorial background in our $B\bar{B}$ MC sample. All four signal channels are combined here. The third and fourth columns show what fraction of the $B$ decays in each group have previously been observed. The fourth column is the product of the second and third, and indicates how much of the estimated combinatorial background is known from other measurements.

| Source                                      | % of total BG (relative) | % $B$ measured (absolute) |
|---------------------------------------------|--------------------------|---------------------------|
| $B \to D^*_+ D^{(*)}$ (+ light hadrons)     |                          |                           |
| ... with $D^*_+ \to \tau\nu$              | 30                       | 90                        | 27            |
| ... with $D^*_+ \to \ell\nu(\phi/\eta/\eta')$ | 10                       | 90                        | 9             |
| $B \to D^{(*)} D^{(*)}$ (+ light hadrons)  | 35                       | 65                        | 25            |
| Both $B \to D^{(*)}\ell\nu$               | 15                       | 100                       | 15            |
| Fake lepton                                 | 5                        | 0                         | 0             |
| Other misreconstructed                      | 5                        | 0                         | 0             |
| Total                                       |                          |                           | 75            |
Simulation discrepancies generally 10% or less in the model. In all figures, the branching-fraction ratio listed in Table VIII, along with relevant branching fractions. In each fit, we reweight events in the simulation. For charge-crossed events (true $B \to D^{(*)} \ell^- \nu$ events where the $B_{\text{tag}}$ and signal $D^{(*)}$ swap a charged particle), the dominant systematic uncertainty is not the branching fraction, but rather the efficiency to reconstruct the $B_{\text{tag}}$ with the wrong charge. We estimate a 10% uncertainty on the modeling of this process, i.e., we generate weights for these events using a Gaussian with a mean of 1 and a width of 0.1. For double-semileptonic events, with both $B$ mesons decaying to $D^{(*)} \ell^- \nu$, again, the dominant uncertainty comes from the probability to misreconstruct a $B_{\text{tag}}$ candidate in this event, and we assume a 10% uncertainty on this number as well. For events in which the signal lepton is misidentified, we assign a 10% uncertainty; the typical fake rate measured in data is 2%-3%, with data-simulation discrepancies generally 10% or less in the momentum ranges of interest. For all remaining sources of combinatorial background, including high-multiplicity $B \to B_{\text{tag}} h (\ldots)$ and $B \to D_{\text{tag}} h (\ldots)$ (where $D$ here represents any charm meson and $h$ any light meson) and other misreconstructed events, we assume a 50% uncertainty in the relevant rates.

In each test, we fit the reweighted MC sample to generate new PDF shapes and recalculate the expected yield of combinatorial events in each channel. Figure 10 shows the effect of this reweighting on the combinatorial BG PDF in the signal channels. We note that the reweighting affects the normalization of the charge-crossed backgrounds but not the shape.

3. $D^{**}$ Modeling

We generate an ensemble of $D^{**}$ models, sampling from the distribution in Table IX. This model is based on the current world averages [16, 28] but includes information from selected recent measurements [33].

| Mode                      | $B$       | MC     |
|---------------------------|-----------|--------|
| $D^+_s \to \tau \nu$      | $(6.4 \pm 1.5) \times 10^{-2}$ | 7.0    |
| $D^+_s \to \eta \ell \nu$ | $(2.5 \pm 0.7) \times 10^{-2}$ | 2.6    |
| $D^+_s \to \eta(\prime) \nu$ | $(8.9 \pm 3.3) \times 10^{-3}$ | 8.9    |
| $D^+_s \to \phi \ell \nu$  | $(2.0 \pm 0.5) \times 10^{-2}$ | 2.0    |
| $D^+_s \to \mu \nu$       | $(5.0 \pm 1.9) \times 10^{-3}$ | 4.6    |
| $B^+ \to D^0 D_s^+$       | $(1.3 \pm 0.4) \times 10^{-2}$ | 1.06   |
| $B^+ \to D^0 D_s^{\ast+}$ | $(9 \pm 4) \times 10^{-3}$    | 9.1    |
| $B^+ \to D^0 D_s^+$       | $(1.2 \pm 0.5) \times 10^{-2}$ | 1.02   |
| $B^+ \to D^0 D_s^{\ast+}$ | $(2.7 \pm 1.0) \times 10^{-3}$ | 2.28   |
| $B^+ \to D^{\ast+} D_s^{\ast+}$ | $(2.7 \pm 1.2) \times 10^{-2}$ | 3.0    |
| $B^+ \to D^{\ast+} D_s^{\ast+}$ | $(5.2 \pm 1.2) \times 10^{-3}$ | 5.2    |
| $B^+ \to D_s^{\ast+} K_s^+$ | $(7.8 \pm 2.6) \times 10^{-3}$ | 7.8    |
| $B^+ \to D_s^{\ast+} \bar{K}_s^+$ | $(1.37 \pm 0.32) \times 10^{-3}$ | 1.9    |
| $B^+ \to D_s^{\ast+} K_s^{0}$ | $(5.3 \pm 1.6) \times 10^{-3}$ | 5.3    |
| $B^+ \to D_s^{\ast+} \bar{K}_s^{0}$ | $(4.7 \pm 1.0) \times 10^{-3}$ | 4.8    |
| $B^+ \to D_s^{\ast+} D_s^{0}$ | $(1.5 \pm 0.4) \times 10^{-3}$ | 0.5    |
| $B^0 \to D^- D_s^+$       | $(8.0 \pm 3.0) \times 10^{-3}$ | 7.4    |
| $B^0 \to D^{\ast+} D_s^+$ | $(1.07 \pm 0.29) \times 10^{-2}$ | 1.03   |
| $B^0 \to D^{-} D_s^{\ast+}$ | $(1.0 \pm 0.5) \times 10^{-2}$ | 0.74   |
| $B^0 \to D^{-} D_s^{\ast+}$ | $(1.9 \pm 0.5) \times 10^{-2}$ | 1.97   |
| $B^0 \to D^{-} D_s^{0}$    | $(1.7 \pm 0.4) \times 10^{-3}$ | 1.7    |
| $B^0 \to D^{-} D_s^{0}$    | $(3.1 \pm 0.6) \times 10^{-3}$ | 3.1    |
| $B^0 \to D^{-} D_s^{0}$    | $(1.18 \pm 0.20) \times 10^{-2}$ | 1.18   |
| $B^0 \to D^{-} D_s^{0}$    | $(6.5 \pm 1.6) \times 10^{-3}$ | 8.1    |
| $B^0 \to D^{-} D_s^{0}$    | $(8.8 \pm 1.9) \times 10^{-3}$ | 8.8    |
| $B^0 \to D^{\ast+} D_s^{-}$ | $(8.1 \pm 1.2) \times 10^{-4}$ | 8.3    |
| $B^0 \to D^{\ast+} D_s^{-}$ | $(10.4 \pm 2.0) \times 10^{-4}$ | 6.7    |
| $B^0 \to D^{\ast+} D_s^{-}$ | $(2.8 \pm 0.7) \times 10^{-4}$ | 2.7    |

FIG. 15: Distributions of fit results for systematic uncertainties due to Monte Carlo statistics, shown for (a) $B^- \to D^{0} \tau^- \nu_{\tau}$, (b) $B^- \to D^{0} \tau^- \nu_{\tau}$, (c) $B^0 \to D^{+} \tau^- \nu_{\tau}$, (d) $B^0 \to D^{0} \tau^- \nu_{\tau}$, (e) $B \to D \tau^- \nu_{\tau}$, and (f) $B \to D^+ \tau^- \nu_{\tau}$. In all figures, the branching-fraction ratio $R$ is shown.

Branching fractions of most of the relevant two-body $B$ decays (and some of the three-body decays as well) have previously been measured. These branching fractions are listed in Table VIII along with relevant branching fractions of the $D^+_s$ meson.

To study systematic uncertainties related to combinatorial background modeling, we perform an ensemble of fits. In each fit, we reweight events in the simulation. For modes listed in Table VIII, we reweight the branching fraction, generating random weights from a Gaussian distribution based on the measured value (for decays involving a $D^+_s$ meson, the weight is the product of weights for both the $B$ and $D^+_s$ decays).
the nonresonant states, maintaining the Monte Carlo ratio of and crossfeed constraints to estimate the systematic un-
parameters given in Section III. We reweight signal and
Caprini, Lellouch, and Neubert [27], with numerical pa-
both cases, we use the form factor parameterization of
FIG. 16: Combinatorial background modeling variation for

| $m_{\text{miss}}^2$ (GeV/c$^2$) | Events / 0.4 (GeV/c$^2$)$^2$ |
|---------------------------|-----------------------------|
|                            |                             |

| $|p_\ell^*|$ (GeV/c) | Events / 60 (MeV/c) |
|-------------------|---------------------|
|                   |                     |

isospin symmetry between charged and neutral $B$ modes.

For each test, we generate random numbers for the six
exclusive modes ($D$, $D^*$, and the resonant $D^{**}$ states),

For each test, we generate random numbers for the six
exclusive modes ($D$, $D^*$, and the resonant $D^{**}$ states),


5. $m_{\text{miss}}^2$ Tail Modeling

Studies in the two kinematic control samples show ac-
ceptable overall agreement between data and simulation for the $m_{\text{miss}}^2$ resolution [see Fig. 8(d)], but suggest that
the simulation may underestimate the ratio of the num-
ber of events in the large $m_{\text{miss}}^2$ tail region to the num-
ber of events near $m_{\text{miss}}^2 = 0$. We estimate that this
tail component of the resolution may be underestimated
by up to 10%. We study systematic effects related to this by reweighting events at large $m_{\text{miss}}^2$, greater than 1 (GeV/c$^2$)$^2$, up by 10%, modifying the PDF shapes for
$B \rightarrow D^*\bar{\tau}\tau$ and $B \rightarrow D^{(*)}\bar{\tau}\tau$. We perform a fit with these modified PDFs and take the difference from the
nominal fit as a systematic uncertainty.

6. $\pi^0$ Efficiency and Crossfeed Constraints

While the systematic uncertainties due to detector ef-
ficiencies (described in more detail in Section IX B 3)
are primarily multiplicative, the efficiencies for $\pi^0$ re-
construction have a large impact on the feed-down ef-
ficiencies and therefore the fit yields. This effect can be
enhanced by the fact that the feed-down constraints are


| Mode | $B_{\text{gen}}$ | $B(D^0)$ | $B(B^-)$ |
|------|----------------|-----------|-----------|
|      | $\mu\sigma$   | $\mu\sigma$ | $\mu\sigma$ |
| $c\ell\nu$ | 10.4 | 10.1$^7$ | 0 | 10.9 | 0 |
| $D$    | 2.10 | 2.14 | 0.14 | 2.29 | 0.16 |
| $D^*$  | 5.6  | 5.54 | 0.25 | 5.94 | 0.24 |
| $D_1$  | 0.56 | 0.47 | 0.08 | 0.58 | 0.06 |
| $D_2$  | 0.37 | 0.35 | 0.07 | 0.46 | 0.08 |
| $D_3$  | 0.20 | 0.46 | 0.09 | 0.45 | 0.09 |
| $D_4$  | 0.37 | 0.85 | 0.20 | 0.83 | 0.20 |
| $D^{*}\pi^0$ | 0.1 | 0.03 | — | 0.029 | — |
| $D\pi^0$ | 0.3 | 0.09 | — | 0.088 | — |
| $D^{*}\pi^\pm$ | 0.2 | 0.06 | — | 0.058 | — |
| $D\pi^\pm$ | 0.6 | 0.18 | — | 0.175 | — |

TABLE IX: $B \rightarrow X_c\ell^-\bar{\nu}_\ell$ branching fractions used in the $D^{**}$ modeling systematic study. The first line, $c\ell\nu$, represents the inclusive semileptonic branching fraction. For the six lines representing the $D$, $D^*$, and $D^{**}$ resonant states, the distribution of these branching fractions is taken to be Gaussian with the given mean and width. For the last four lines, representing the nonresonant $D^{**}$ states, the ranges of variation are not shown in this table; their distribution is determined by the inclusive rate and the other exclusive modes, as described in the text. The generated branching fractions, $B_{\text{gen}}$, are the same for charged and neutral $B$ mesons. All numbers are in %.

4. $B \rightarrow D^{(*)}$ Form Factors

We reweight the form factors of both signal $B \rightarrow D^{(*)}\tau^-\bar{\nu}_\tau$ and normalization $B \rightarrow D^{(*)}\bar{\tau}\tau$ decays. In both cases, we use the form factor parameterization of Capri\"u, Lellouch, and Neubert [27], with numerical pa-
parameters given in Section III. We reweight signal and
defined as the ratio of the number of events reconstructed in the \( D\ell^- \) channel to that in the \( D^*\ell^- \) channel, which move in opposite directions as the \( \pi^0 \) efficiency is varied.

We generate an ensemble of fits by varying the \( \pi^0 \) efficiency within its uncertainty, 3.0\% per \( \pi^0 \). The resulting changes in the feed-down constraints for both signal and background modes are propagated through the signal fit to estimate the resulting systematic uncertainties.

7. \( D^{**}\ell^-\nu_\ell \) Feed-down

We assign an additional systematic uncertainty on \( D^{**}\ell^-\nu_\ell \) feed-down rates due to the fact that the \( \pi^0 \) mesons involved in feed-down processes typically have low momentum, while the 3.0\% systematic uncertainty mentioned above is derived from a control sample with a broad spectrum. Since we float the constraints describing \( D^* \to D \) feed-down in the fit, \( D^* \) feed-down processes are insensitive to systematic effects due to the \( \pi^0 \) efficiency at low momentum. The \( D^{**} \) feed-down constraints, however, are taken from simulation and can therefore be affected.

We compare the fitted values of the \( D^* \to D \) feed-down rates to the simulation to estimate that the efficiency for low-momentum \( \pi^0 \) mesons is correctly modeled to within 10\%. We generate an ensemble of fits in which we vary the \( \pi^0 \) reconstruction efficiency \( \pm 10\% \) for \( \pi^0 \) mesons with momentum less than 300 MeV/c. We generate new PDFs and feed-down constraints which we propagate through the signal fit to estimate the systematic uncertainties.

8. \( B \to D^{**}\tau^-\nu_\tau \) Abundance

We vary the fraction of \( B \to D^{**}\tau^-\nu_\tau \) events in the \( D^{**} \) samples by generating random numbers from a Gaussian distribution with mean 1.0 and width 0.3, equivalent to a \( \pm 30\% \) variation. For each test, we generate new PDFs and crossfeed constraints to estimate the systematic uncertainty.

B. Multiplicative Systematic Uncertainties

1. Monte Carlo Statistics

The dominant multiplicative systematic uncertainty is due to limited Monte Carlo statistics. The various MC samples are independent of one another, so that there is no cancellation between the signal and normalization.

2. Bremsstrahlung and Final-State Radiation

Based on a control sample of identified electrons and studies in MC samples, we estimate the uncertainty on reconstruction efficiency due to Bremsstrahlung and final-state radiation effects to be 2.1\%. This uncertainty applies to both signal and normalization modes, however, and so the effect on the relative efficiency is expected to cancel. The fractions of events in which a photon is radiated are nearly the same between signal and normalization modes, within statistical precision of 10\%; we therefore treat the uncertainty between signal and normalization modes as 90\% correlated to calculate the final systematic uncertainty.

3. Detector Efficiencies

We estimate systematic uncertainties related to the detector efficiencies—track and neutral reconstruction and charged particle identification—by studying these efficiencies in several control samples in both data and simulation. We correct the MC efficiencies to match those seen in the data, and we take the statistical precision of these studies as an estimate of the systematic uncertainty on absolute efficiencies.

Since we normalize our signals to \( B \to D^{(*)}\ell^-\nu_\ell \), we calculate systematic uncertainties on the relative efficiency, treating uncertainties on the signal and normalization modes as correlated. The degree of correlation, and therefore, the degree to which the uncertainty cancels, is determined by the kinematics of the two samples. For most of the final state particles, the kinematic distributions are very similar between signal and normalization modes and so the systematic uncertainty cancels almost entirely. For the charged leptons, however, the momentum spectra are very different between signal and normalization (see Fig. 2), and so the associated systematic uncertainty is larger.

4. Hadronic Daughter Branching Fractions

We reconstruct both signal and normalization modes using the same set of final states, so uncertainties due to the branching fractions of these states very nearly cancel. (The \( D^{(*)} \) momentum spectra are slightly different between signal and normalization modes, so this cancellation is not perfect.) We take the uncertainty on each of the reconstructed \( D^*, D, K_0^0 \), and \( \pi^0 \) decay modes from [16] and propagate each of these uncertainties through to the relative efficiency, using the relative abundance of each decay chain in the signal and normalization MC samples to determine the correlation and the degree of cancellation.

5. Leptonic \( \tau \) Branching Fraction

The \( \tau \) branching fraction \( \mathcal{B}(\tau^- \to \ell^-\nu_\ell\nu_\tau) \) appears only in the denominator of Eq. [18] and therefore contributes a 0.2\% systematic uncertainty on all modes [16] without cancellation.
X. RESULTS

Table X summarizes the results from two fits, one in which all four signal yields can vary independently, and the second $B^−\to\overline{B}^0$ constrained fit with $R(D^+) = R(D^0)$ and $R(D^{+*}) = R(D^{*0})$. We observe approximately 67 $B \to D\tau\nu_\tau$ and 101 $B \to D^{+}\tau\nu_\tau$ signal events in this $B^−\to\overline{B}^0$-constrained fit, corresponding to signal branching-fraction ratios of $R(D) = (41.6 \pm 11.7 \pm 5.2)\%$ and $R(D^*) = (29.7 \pm 5.6 \pm 1.8)\%$, where the first error is statistical and the second systematic. Normalizing these to known $\overline{B}^0$ branching fractions, we obtain $B(B \to D\tau\nu_\tau) = (0.86 \pm 0.24 \pm 0.11 \pm 0.06)\%$ and $B(B \to D^{+}\tau\nu_\tau) = (1.62 \pm 0.31 \pm 0.10 \pm 0.05)\%$, where the third error is from that on the normalization branching fraction.

Table X also gives the significances of the signal yields. The statistical significance is determined from $\sqrt{2\Delta(\ln L)}$, where $\Delta(\ln L)$ is the change in log-likelihood between the nominal fit and the no-signal hypothesis. The total significances are determined by including the systematic uncertainties on the fit yields in quadrature with the statistical errors. In the $B^−\to\overline{B}^0$-constrained fit, the signal significances are $3.6\sigma$ and $6.2\sigma$ for $R(D)$ and $R(D^*)$, respectively.

The statistical correlation between $R(D)$ and $R(D^*)$ is $-0.51$ in the $B^−\to\overline{B}^0$-constrained fit. This correlation is due to the fact that most of the events at large $m_{\text{miss}}^2$ are either $B \to D\tau\nu_\tau$ or $B \to D^{+}\tau\nu_\tau$ signal events, and increasing either of the two signal yields in the fit necessarily decreases the other. The systematic uncertainties have a correlation of $-0.03$ between $R(D)$ and $R(D^*)$; most of the systematic uncertainties have large negative correlations for the same reason that the statistical uncertainty does, but the combinatorial background uncertainty affects both signal yields in a coherent manner and so contributes a large positive correlation. The sum of the two branching fractions, taking all correlations into account, is $B(B \to D^{(+)}\tau\nu_\tau) = (2.48\pm0.28\pm0.15\pm0.08)\%$.

Figures 17 and 18 show the observed $q^2$ distributions in the four signal channels in the low and high $m_{\text{miss}}^2$ regions, respectively. The histograms in these figures are taken from MC samples of the various components, with each component scaled to match the yield in the $B^−\to\overline{B}^0$-constrained fit; since $q^2$ is not a fit variable, we cannot show a projection of a continuous PDF as was done in Figs. 10–14. As before, we observe good agreement between the data and the expectation from simulation, in both the low and high $m_{\text{miss}}^2$ regions. Since the $q^2$ distribution is highly dependent on the form factor model, we note that the CLN model describes both normalization and signal events within the available statistics.

Table XI summarizes the results of several crosschecks, including splitting up the sample according to lepton flavor, lepton charge, and data-taking period. We have done these checks by performing “cut-and-count” analyses, both in the data and in simulated event samples. In all cases, the results in data are consistent with our expectations from simulation. The first row in this table shows the fraction of events with muon candidates in data and simulation, both for the full event sample and for the signal-sensitive region in $m_{\text{miss}}^2$. Electron identification is more efficient than muon ID, which is why the muon fraction in the final sample is less than

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8 We use $[18]$ to normalize the four individual branching fractions. For the $B^−\to\overline{B}^0$-constrained measurement, we use our own averages of the values in $[10]$: $B(\overline{B}^0 \to D^{+}\tau\nu_\tau) = (2.07 \pm 0.14)\%$ and $B(\overline{B}^0 \to D^{+}\tau\nu_\tau) = (5.46 \pm 0.18)\%$. 

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FIG. 17: (Color online) $q^2$ distributions of events in the four final states $D^{*0}\ell^+$, $D^{0}\ell^−$, $D^{+}\ell^−$, and $D^{*+}\ell^−$, shown in the normalization region, $m_{\text{miss}}^2 < 1$ (GeV/c$^2$)$^2$. The data are shown as points with error bars. The shaded histograms are taken from MC samples with normalizations from the fit to data. The components are shaded as in Fig. 10.

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TABLE XI: Crosscheck studies, splitting the data according to lepton flavor, lepton charge, and running period. The first row shows the fraction of events with muon candidates for both data and MC samples, for both the full event sample and for the signal-sensitive region $m_{\text{miss}}^{2} > 1$ (GeV/$c^2$)$^2$. The second row shows fractions of events with positively charged lepton candidates, and the third row shows the fractions of events recorded in Run 4. In all cases, the data are consistent with the simulation and with expectations.

| Mode       | $\mathcal{N}_{\text{sig}}$ | $\mathcal{N}_{\text{norm}}$ | $(\Delta R/R)_{\text{fit}}$ | $(\Delta R/R)_{R}$ | $R$ | $B$ | $\sigma_{\text{tot}}$ |
|------------|-----------------------------|-----------------------------|----------------------------|---------------------|-----|-----|-----------------|
| $B^- \rightarrow D^0 \tau^- \nu_\tau$ | 35.6±19.4                   | 347.9±23.1                  | 15.5                       | 1.6                 | 31.4±17.0±4.9 | 0.67±0.37±0.11±0.07 | 1.8 (1.8) |
| $B^- \rightarrow D^{0*} \tau^- \nu_\tau$ | 92.2±19.6                   | 1629.9±63.6                 | 9.7                        | 1.5                 | 34.6±7.3±3.4 | 2.25±0.48±0.22±0.17 | 5.3 (5.8) |
| $\bar{B}^0 \rightarrow D^+ \tau^- \nu_\tau$ | 23.3±7.8                    | 150.2±13.3                  | 13.9                       | 1.8                 | 48.9±16.5±6.9 | 1.04±0.35±0.15±0.10 | 3.3 (3.6) |
| $\bar{B}^0 \rightarrow D^{+*} \tau^- \nu_\tau$ | 15.5±7.2                    | 482.3±25.5                  | 3.6                        | 1.4                 | 20.7±9.5±0.8 | 1.11±0.51±0.04±0.04 | 2.7 (2.7) |
| $B^- \rightarrow D^- \tau^- \nu_\tau$ | 66.9±18.9                   | 497.8±26.4                  | 12.4                       | 1.4                 | 41.6±11.7±5.2 | 0.86±0.24±0.11±0.06 | 3.6 (4.0) |
| $B^- \rightarrow D^+ \tau^- \nu_\tau$ | 101.4±19.1                  | 2111.5±68.1                 | 5.8                        | 1.3                 | 29.7±5.6±1.8 | 1.62±0.31±0.10±0.05 | 6.2 (6.5) |

The statistical correlation between $R(D)$ and $R(D^*)$ in this fit is $-0.51$.

XI. CONCLUSIONS

We have presented measurements of the branching fractions for the decays $B \rightarrow D^{(*)} \tau^- \nu_\tau$ and $B \rightarrow D^{(*)} \tau^- \nu_\tau$, determined relative to the corresponding decays to light leptons. We measure the branching-fraction ratios for four individual $D^{(*)}$ states, as well as two $B^- - \bar{B}^0$-constrained ratios

$$R(D^0) = (31.4 \pm 17.0 \pm 4.9)\%$$
$$R(D^{*0}) = (34.6 \pm 7.3 \pm 3.4)\%$$
$$R(D^+) = (48.9 \pm 16.5 \pm 6.9)\%$$
$$R(D^{*+}) = (20.7 \pm 9.5 \pm 0.8)\%$$
$$R(D) = (41.6 \pm 11.7 \pm 5.2)\%$$
$$R(D^*) = (29.7 \pm 5.6 \pm 1.8)\%$$

where the first uncertainty is statistical and the second is systematic. The significances of these signals are $1.8\sigma$, $5.3\sigma$, $3.3\sigma$, $2.7\sigma$, $3.6\sigma$, and $6.2\sigma$, respectively. The statistical and systematic uncertainties on $R(D)$ and $R(D^*)$ have correlations of $-0.51$ and $-0.03$, respectively.

From these branching-fraction ratios and known branching fractions of the normalization modes $B \rightarrow D^{(*)} \ell^- \nu_\ell$, we derive the absolute branching fractions...
FIG. 18: (Color online) $q^2$ distributions of events in the four final states $D^{\ast\ell^-}$, $D^0\ell^-$, $D^{\ast+}\ell^-$, and $D^{+}\ell^-$, shown in the signal region, $m_{\text{miss}} > 1 \text{ (GeV/c}^2\text{)}$. The data are shown as points with error bars. The shaded histograms are taken from MC samples with normalizations from the fit to data. The components are shaded as in Fig. 10.

\begin{align*}
B(B^- \to D^0\tau^-\tau_\tau) & = (0.67 \pm 0.37 \pm 0.11 \pm 0.07)\% \\
B(B^- \to D^{\ast0}\tau^-\tau_\tau) & = (2.25 \pm 0.48 \pm 0.22 \pm 0.17)\% \\
B(B^\ast \to D^{\ast+}\tau^-\tau_\tau) & = (1.04 \pm 0.35 \pm 0.15 \pm 0.10)\% \\
B(B^\ast \to D^{\ast+}\tau^-\tau_\tau) & = (1.11 \pm 0.51 \pm 0.04 \pm 0.04)\% \\
B(B \to D\tau^-\tau_\tau) & = (0.86 \pm 0.24 \pm 0.11 \pm 0.06)\% \\
B(B \to D^{\ast}\tau^-\tau_\tau) & = (1.62 \pm 0.31 \pm 0.10 \pm 0.05)\% ,
\end{align*}

where the third uncertainty reflects that of the normalization mode branching fraction.

The measurement of $B(B^\ast \to D^{\ast+}\tau^-\tau_\tau)$ is consistent with the Belle result [17]. The branching-fraction ratios $R(D)$ and $R(D^\ast)$ are about 1σ higher than the SM predictions but, given the uncertainties, there is still room for a sizeable non-SM contribution.

We have also presented distributions of the lepton momentum $|p^*_\ell|$ and the squared momentum transfer $q^2$ for $B \to D^{(\ast)}\tau^-\tau_\tau$ events. In all cases, these distributions are consistent with expectations based on the SM and the CLN form factor model with measured form factors.

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