Optical Dynamics at the Au/ZnPc Interfaces

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In this work, the optical dynamics and the structural properties of the zinc phthalocyanine which are coated onto 150 nm thick Au substrates are studied by the X-ray diffraction and optical spectrophotometry techniques. The Au/ZnPc interfaces appear to be strongly affected by the large lattice mismatches at the interface. It is observed that the coating ZnPc onto Au substrates increases the light absorbability by 4.7 and 128.2 times in the visible and infrared regions of light, respectively. Au substrates activated the free carrier absorption mechanism in the ZnPc thin films in the infrared range of light. In addition, the transparent Au substrates forced narrowing the energy band gap in both of the Q and B bands. It also increased the dielectric constant value by ~3.5 times in the IR range. The enhancements in the optical properties of ZnPc that resulted from the thin Au substrates make the ZnPc more suitable for optoelectronic, nonlinear optical applications and for electromagnetic energy storage in the infrared range of light.

Keyword: Au/ZnPc; X-ray diffraction; optical absorption; band gap; dielectric dispersion.

1. Introduction

Zinc phthalocyanine (ZnPc) is a compound that exhibits smart features nominating it for many applications. The kind of application of this material depends on the substrates they grow onto. As for examples, ZnPc thin films which are coated onto Fe3O4 are reported to be beneficial or photodynamic therapy and for magnetic resonance imaging1. ZnPc/carbon nitride interfaces are observed to display high efficiency in photoelectrocatalytic reduction of CO2 under visible light. In addition, heterojunctions composed of C60 and zinc phthalocyanine which are used to fabricate organic photodiodes displayed a photoresponsivity of 5.21 A/ W when irradiated with light of intensity of 68 µW/cm2. Moreover, as a nonlinear optical media, when tested with the Z-scan technique using 532 nm pulsed lasers, the metal phthalocyanine/graphene oxide (GO) interfaces displayed third order nonlinear optical absorption1. Furthermore, when the CdS/CdSe solar cells were coated with ZnPc the short current density, open circuit voltage and photocell efficiency increased to 11.3 mA/cm2, 0.66 V and 4.01%, respectively2.

Literature data mentioned the novelty of ZnPc for use in the technology of nanogenerator3-7 which are well known for their high energy conversion efficiency7. Recently, Zinc phthalocyanine nanorods which are grown onto flexible aluminum foil (Al) are nominated for use in the nanogenerators technology. Al play the role of substrate to grow nanorods and as bottom electrodes. For such system an open circuit voltage of 0.524 mV for ZnPc rods of lengths of 100 nm is achieved. This open circuit voltage reaches 0.968 mV when the length of the rod is increased to 300 nm8.

The wide range of applications and the effect of the type of substrate on the performance of ZnPc thin films, which are additionally summarized in Table 1, motivated us to study the relation between transparent Au layers and ZnPc thin films. This type of interfaces may open the doors for new class of applications that need electrical contacts and high optical transparency to carryout dual operations. For this reason, here in this work, we will investigate the structural and optical properties of ZnPc films which are coated onto glass substrates and compare them with those coated onto 150 nm thick Au substrates. The feature of the transparent Au substrate lay in its high conductivity and high light transparency. The effect of Au on the crystal structure, optical absorption, light absorbability, energy band gap and dielectric constant will be reported and discussed. Enhancement of these physical parameters will make the ZnPc more attractive for use in optoelectronics as multifunctional devices.

2. Experimental details

Both of the Au and ZnPc thin films are coated onto ultrasonically cleaned glass substrates with the help of VCM-600 thermal evaporator under vacuum pressure of 10-5 mbar. The deposition parameters during the evaporation cycle are listed in Table 2. The respective thicknesses of the films were 150 and 600 nm. The evaporation sources were Au wires and ZnPc nano-powders. The Au substrates were also employed to grow the 600 nm thick layers of ZnPc. The thicknesses of the films were measured with Inficon STM-2 thickness monitor which have resolution of 0.037 Å. The measurement intervals of the monitor are 0.1 s spaced. The monitor detecting face was set parallel to the glass substrates so that it read the same thicknesses of the grown films. As confirmative procedure, the produced

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Table 1: Optoelectronic devices comprising ZnPc in its structure.

| Device structure | Class of applications | Reference |
|------------------|-----------------------|-----------|
| Fe₂O₃/ZnPc       | Photodynamic therapy and for magnetic resonance imaging | Qin et al.¹ |
| ZnPc/C₆H₆   | Photoelectrolytic reduction of CO₂ | Zheng et al.² |
| C₆₀/ZnPc        | Organic photodiodes with high responsivity | Ding et al.³ |
| ZnPc/C₆H₆O₂*   | Third order nonlinear optical absorption | Zu et al.⁴ |
| CdS/CdSe/ZnPc   | Solar Cells            | Chen et al.⁵ |
| Al/ZnPc         | Nanogenerators         | Godfrey et al.⁶ |

*Famous name: Graphene oxide

Table 2: The deposition parameters for ZnPc films.

| Parameter                             | Value           |
|---------------------------------------|-----------------|
| Pre-deposition vacuum pressure (mbar) | 4×10⁻⁵          |
| Evaporation Boat type                 | Tungsten (0.1mm thick) |
| Weight of ZnPc powder (gr)            | 0.14            |
| Target to substrate distance (cm)     | 10              |
| Temperature of substrate holder (°C)  | 25              |
| Evaporation time (min.)               | 13              |
| Evaporation rate (Å/s)                | 8.4             |

films were measured with the help of surface roughness tester-profilometer (Model TR-200 plus). The profilometer reads 655 nm for the ZnPc films which differ by 9.1% from the monitor. The arithmetic average roughness (R_a) was found to be 55 nm. The thickness of the Au layer being 150 nm was optimized through successive cycles of evaporations to determine the best position of the glass substrates above the evaporation source and then confirming the thicknesses by the profilometer which is capable of recording thicknesses down to 5.0 nm. The grown thin films of Au/ZnPc are structurally and optically characterized using MiniFlex-600 X-ray diffraction unit with scanning speed of 0.5°/min and Thermoscientific Evolution 300 spectrophotometer equipped with VEE MAX II reflectometer. The optical measurements were carried out in the spectral range of 300-1100 nm.

It is worth mentioning that the thickness of the samples being 600 nm is selected for the preparation of the Au/ZnPc interfaces after various trials to optimize stable crystalline structure. Films of thicknesses in the range of 50-500 nm did not display any experimental intensive peaks in the XRD patterns. While, on the other hand, films of thicknesses of 600-1000 nm exhibited polycrystalline phase. Thus, the minimum thickness that reveal crystalline phase and highest transmittance of the crystalline phase was at 600 nm.

3. Results and Discussion

The X-ray diffraction (XRD) patterns for zinc phthalocyanine thin films which are coated onto glass and transparent Au substrates are shown in Figure 1. As mentioned, the transparent Au substrate thickness (d) was 150 nm. It is clear from the figure that the replacement of glass by Au has significant effect on the intensity and diffraction angle (2θ) of the maximum peak. Particularly, as appears in inset-1 of Figure 1, the peak which was observed at 2θ = 7.32° and exhibit intensity 2.66×10⁵ a.u shift to 2θ = 7.25° and display intensity value of 1.34×10⁶ a.u. The maximum peak of the XRD patterns which is usually assigned to the monoclinic phase of ZnPc with unit cell parameters of \( a = 26.333 \, \text{Å}, \, b = 3.8133 \, \text{Å}, \, c = 23.7118 \, \text{Å}, \) and \( \beta = 94.2845° \) and space group: \( C_{2h} \) are oriented along the (002) direction (JCPDS card number 21-1985). As also observable from the XRD patterns of the Au/ZnPc films, additional peak that is assignable to the face centered cubic gold with plane orientation along the (111) direction is presented. The calculated lattice parameter of gold is 4.05 Å(JCPDS Card No. 04-0784). As both of the ZnPc peaks exhibit the same broadening value (\( \beta = 0.375° \)) no effect of Au substrates on the crystallite size or microstrain is expected. However, the large differences between the lattice parameters of the monoclinic unit cells of ZnPc and cubic Au turn the attention to the large lattice mismatch between the Au and ZnPc layers. Large lattice mismatch is mentioned leading to the formation of a high density of structural defects like threading dislocations which strongly affect the physical properties of thin films⁷. Large lattice mismatches are also reported to cause small binding force. In this case, the ZnPc thin film exhibits spontaneous nucleation crystallization⁸. The spontaneous nucleation process could be accounted as a main reason for the shift in the diffraction angle position and for the lower intensity of the XRD patterns in the presence of Au substrates. Thus, in the light of these analyses, the decrease in the intensity of the XRD patterns upon replacement of substrates can be assigned to the decrease in the number of oriented planes along the (002) direction as a result of large lattice mismatches and small binding forces. Earlier
studied on gold-phthalocyanine interactions indicated that, in general, the different growth modes of phthalocyanine molecules onto metal substrates allow direct control of the \( b \)-axis or the one-dimensional chain direction in thin films\(^{11} \). The molecule-molecule and molecule-substrate interactions end with variations of the atomic ordering on the substrate. The stronger interactions between gold-phthalocyanine as compared to the phthalocyanine-phthalocyanine interaction, makes the first layer of phthalocyanine grows in a flat layer on the surface and the subsequent layers follow the first layer.

In an attempt to observe the effect of the Au layer on the surface morphology of ZnPc, the scanning electron microscopy technique was employed. Optimizing clear images from the surface of the ZnPc samples which are grown onto glass was not possible. However, for the Au/ZnPc films the scanning electron microscopy revealed the image which is shown In inset-2 of Figure 1. The inset which represents an enlargement of 10,000 times displayed randomly distributed grains. On average, the grains are irregularly shaped exhibited sizes of 169 nm. While the random distribution of grains means the weak crystallinity (as also suggested by XRD) of the material, this large size of grains indicates the accumulation of many crystallites to form one grain. Generally, the random plane orientations of the accumulated crystallites to form a grain decrease the degree of crystallization.

Figure 2(a) display the absorption coefficient (\( \alpha \)) spectra for the ZnPc thin films which are coated onto glass and transparent Au substrates. The absorption coefficient is calculated from the measured transmittance (\( T \)) and reflectance (\( R \)) using the relation, \( \alpha = \frac{\ln (1 - R_{av})}{d} \) (\( R_{av} = \frac{R}{\cos^2 \theta} \)). It is evident from Figure 2(a) that the absorption coefficient spectra of the ZnPc thin films are remarkably influenced by the Au substrate in the incident photon energy range of 3.30-4.10 eV (376-300 nm) and in the infrared range (IR) of light (1.55-1.14 eV, 800-1100 nm ) as well. The IR region for the Au/ZnPc is shown by green colored circles and enlarged in the inset of Figure 2(b). In this region of light, while the glass/ZnPc show near zero values of absorption coefficient, the spectra of the Au/ZnPc increases with decreasing photon energy. Such phenomena are known as free carrier absorption and is ascribed to the weighted contribution of both acoustical phonon scattering and longitudinal optical phonon scattering\(^{12} \). Similar to our observation, the enhanced infrared absorption in CdMnTe is mentioned to be caused by the lattice absorption and the free carrier absorption mechanisms. The lattice absorption is believed to arise from the lattice distortion in CdMnTe crystal which is caused by dislocations, inclusions and impurities\(^{13} \). The carrier movement which is affected by phonon scattering transfers the energy to lattice when irradiated by IR light leading to the free carrier absorption. The effect of the Au substrates on the light absorbability of ZnPc is easily observable from the absorption coefficient ratios (\( R_{av} \)) of the Au/ZnPc to glass/ZnPc. \( R_{av} \) spectra are shown in Figure 2(b).

In general, the light absorbability are larger than one in most of the spectral range. However, it shows large increase and exhibit one local maxima of 4.70 at 2.72 eV (457 nm) and one absolute maxima of 128.2 at 1.29 eV (963 nm). This is a relatively high value indicating the better performance of the Au/ZnPc over glass/ZnPc films. The large increase in the light absorbability make the Au/ZnPc films promising materials as IR sensors and attractive for use in solar cells\(^{14} \). The enhancement of light absorption upon replacement of glass by Au is mainly attributed to the free carrier absorption which is motivated by the Au substrates in the IR range.

On the other hand, the employment of Tauc’s equation \( \alpha \propto E^{2} \) to determine the effects of transparent Au substrates on the optical energy band gaps revealed the straight lines which are shown in Figure 3(a) and 3(b). While the \( E \)-axis crossings in the Q –band of ZnPc which is shown in Figure 3(a) display a shift in the value of the energy band gap from 1.680 to 1.663 eV, the \( E \)-axis crossing in the B-band display a shift from 3.175 to 3.095 eV upon replacement of glass by Au substrates. The shrinkage in the energy band gap of ZnPc which results from Au substrates can be ascribed to the stronger interactions between the Au and phthalocyanine layers compared to the interactions between phthalocyanine - phthalocyanine layers\(^{11} \). It also could be accounted to the quantum size effect, its stated that\(^{17} \) when the particle size is in the nanoscale regime a higher energy shift in the conduction band happens owing to the larger surface-to-volume ratio. Because of this ratio majority of the constituent atoms can reside on the particle surface, and induces the efficient transfer of charge carriers to suitable regions on the surface. The narrowing of the
energy band gaps in both of the Q and B bands could also be assigned to the orbital overlapping\textsuperscript{15} between the Au ($4f^{14}5d^{10}6s^1$) orbitals and Zn ($4d^{10}4s^2$) atoms in ZnPc. The Au metal can add one more free electron to the ZnPc orbitals as it reaches $6s^1$ states. In general, the optical absorption in the Q and B bands in ZnPc originates from the highest occupied molecular orbitals (HOMO) to the lowest unoccupied molecular orbitals (LUMO) states. The ground state of the HOMO and SHOMO (singlet highest occupied molecular orbital) is due to the $a_{1g}(\pi)$ and $a_{2g}(\pi^*)$ orbitals. The first excited LUMO level is $e_g(\pi^*)$\textsuperscript{18,19}. The electronic transitions from $a_{1g}(\pi)$ to $e_g(\pi^*)$ result in Q- and B-band absorptions near 2.0 and 3.65 eV. The Q-band splits into two bands centered at 1.75 and 2.0 eV because of the molecular vibration limitations in phthalocyanine microcycles\textsuperscript{18,19}. Participation of the Au ($6s^1$)orbitals increases the number of electrons that are available to support these optical transitions which could also account for the changes in the energy band gap values.

Figure 4 illustrates the real part ($\varepsilon_r$) of the dielectric spectra for the glass/ZnPc and Au/ZnPc thin films. The dielectric spectra are calculated with the help of Fresnel’s equations using the measured $T$ and $R$ data\textsuperscript{18,16}. The figure displays broad peaks of $\varepsilon_r$ spectra at critical energy values of 3.14, 2.68, 1.94 and 1.62 eV for glass/ZnPc. These peaks appear at 2.77, 1.94, 1.72 and 1.25 eV for the Au/ZnPc films. While the peaks at 3.14 eV and 2.66 eV are assigned to the transitions in the B bands, the critical energies at 1.94 and 1.62 eV in glass/ZnPc are assigned to the Q-band and SHUMO transitions in ZnPc. The shifted peaks which appeared as a result of Au substrates may be assigned to the same reasons we mentioned for the attenuations in the energy band gap. The most import point that is worthy of consideration in the dielectric spectra is the large enhancement in the value of the dielectric constant upon coating of ZnPc onto Au substrates. The inset shows the ratio of the dielectric constants. In the current study we have shown that interfacing of zinc phthalocyanine with Au thin layers of thicknesses of 150 nm could significantly improve the optical properties of ZnPc and make it more appropriate for optoelectronic technology applications. Both of the light absorbability and dielectric constant are remarkably increased when the films were coated onto Au substrates. It is also observed that the Au caused slight shrinkage in the energy band gaps in both of the B and Q bands of ZnPc. The large values of the dielectric constant at 1.30 eV makes the Au/ZnPc films attractive for electromagnetic energy storage.

**4. Conclusions**

In the current study we have shown that interfacing of zinc phthalocyanine with Au thin layers of thicknesses of 150 nm could significantly improve the optical properties of ZnPc and make it more appropriate for optoelectronic technology applications. Both of the light absorbability and dielectric constant are remarkably increased when the films were coated onto Au substrates. It is also observed that the Au caused slight shrinkage in the energy band gaps in both of the B and Q bands of ZnPc. The large values of the dielectric constant at 1.30 eV makes the Au/ZnPc films attractive for electromagnetic energy storage.

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