Bimolecular reactions of CH$_2$CN$^{2+}$ with Ar, N$_2$ and CO: reactivity and dynamics†

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The reactivity, energetics and dynamics of bimolecular reactions between CH$_2$CN$^{2+}$ and three neutral species (Ar, N$_2$ and CO) have been studied using a position sensitive coincidence methodology at centre-of-mass collision energies of 4.3–5.0 eV. This is the first study of bimolecular reactions involving CH$_2$CN$^{2+}$, a species relevant to the ionospheres of planets and satellites, including Titan. All of the collision systems investigated display two collision-induced dissociation (CID) channels, resulting in the formation of C$^+$ + CH$_2$N$^+$ and H$^+$ + HC$_2$N$. Evidence for channels involving further dissociation of the CID product HC$_2$N$^+$, forming H + CCN$, were detected in the N$_2$ and CO systems. Proton-transfer from the dication to the neutral species occurs in all three of the systems via a direct mechanism. Additionally, there are product channels resulting from single electron transfer following collisions of CH$_2$CN$^{2+}$ with both N$_2$ and CO, but interestingly no electron transfer following collisions with Ar. Electronic structure calculations of the lowest energy electronic states of CH$_2$CN$^{2+}$ reveal six local geometric minima: both doublet and quartet spin states for cyclic, linear (CH$_3$CN), and linear isocyanide (CH$_2$NC) molecular geometries. The lowest energy electronic state was determined to be the doublet state of the cyclic dication. The ready generation of C$^+$ ions by collision-induced dissociation suggests that the cyclic or linear isocyanide dication geometries are present in the [CH$_2$CN]$^{2+}$ beam.

1. Introduction

Doubly-charged positive ions (dications) can be found in a variety of energised environments including the ionospheres of planets and their satellites. A number of atomic dications have been detected in planetary ionospheres, and, recently, the molecular dication CO$_2^{2+}$ has been detected in the ionosphere of Mars; the first detection of a molecular dication in such an environment. The presence of molecular dications in planetary ionospheres is predicted by modelling, and the CO$_2^{2+}$ dication has also been detected in a cometary gas envelope. Despite, in most cases, their inherent thermodynamic instability towards charge separation, the metastable electronic states of small molecular dications have been shown to possess lifetimes sufficient to allow collisions with other species in ionospheric and other environments. Indeed, the lifetimes of atomic dications in planetary ionospheres are expected to be primarily determined by collisional processes. Significant bimolecular reactivity has been observed following collisions of both atomic and molecular dications with neutral species. This bimolecular reactivity suggests that dication chemistry can play a role in ionospheric processes; for example, dications could be involved in the chemistry of complex molecule formation through carbon chain-growth reactions. In order to identify dication reactions of ionospheric interest, experiments to probe dicaticonic reactivity are vital. For example, recent experimental work has identified the role molecular dications can play in atmospheric erosion processes, such as that involving CO$_2^{2+}$ dissociation in the atmosphere of Mars. This paper presents an investigation of the reactions between the molecular dication CH$_3$CN$^{2+}$ and Ar, N$_2$ and CO. This work provides the first probe of the reactivity, reaction mechanisms and energetics of this dication, enhancing our understanding of the role of such species in energised environments such as planetary atmospheres.

The atmosphere of Titan is primarily comprised of molecular nitrogen (N$_2$, $\sim$95%) and methane (CH$_4$, $\sim$5%). It follows that nitriles, organic compounds containing a $-CN$ moiety, are also relevant to the atmospheric chemistry of this satellite. Indeed, a number of nitrile compounds, including acetonitrile, have been detected in the atmosphere of Titan. Acetonitrile, CH$_3$CN, is the simplest nitrile, and has also been detected in the atmosphere of the Earth, as well as in the interstellar medium (ISM). A recent investigation of the electron ionization and photoionization of CH$_3$CN detected the formation of long-lived molecular dications including, predominantly, CH$_3$CN$^{2+}$. In ionospheres, dications are

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readily formed from their precursors by solar radiation and by collisions with high-energy particles.14 Therefore, the CH$_2$CN$^{2+}$ dication is likely to be produced from CH$_3$CN in the ionosphere of Titan, as well as in the atmospheres of other planets and satellites. In addition to its relevance in planetary ionospheres, the neutral radical CH$_3$CN has been detected in the interstellar medium, where nitriles are considered to play an important role in mechanisms for formation of complex organic molecules.36,39–41

Our choice of neutral collision targets is stimulated by their simplicity and relevance to planetary atmospheres. Argon is often present in the atmospheres of planets and satellites, for example in the atmospheres of Earth, Mars and Titan.30,44–49 Molecular nitrogen is the dominant species in the atmosphere of the Earth, but also in the atmosphere of Titan.13,19,45–48,50–52 While CO has been detected in comets,53–55 is widely distributed across the ISM;56,57 and is also present in the atmospheres of various planets and satellites including the Earth,58 Venus,59 and Titan.60–62

To the authors’ knowledge, no bimolecular reactions involving CH$_2$CN$^{2+}$ have previously been reported. However, comparisons can be made with the reactivity of similar molecular dications when they encounter the neutral collision partners used in this study. The reactions resulting from dication collisions with Ar have been well studied.63–76 Along with other channels, proton-transfer (PT) forming ArH$^+$ was detected in collisions between Ar and the hydrogen-containing dications CHCl$^{2+}$ and C$_2$H$_2$Cl$^{2+}$.77,78 Additionally, Ar + CHCl$^{2+}$ collisions generated HCCCl$^+$.78 The reactions resulting from dication collisions with N$_2$ have been the subject of several previous studies, including recent work by the current authors.63,69–71,79–86 PT reactions have been observed following the interactions of N$_2$ with C$_2$H$_2^+$, resulting in the formation of N–H bonds.87 CO has been used as a collision partner in several studies of dication reactivity.70,84,88–91 Collisions between HCl$^{2+}$ and CO resulted in PT forming COH$^+$.87 In the light of these previous studies, we may hypothesise that PT is a likely consequence of CH$_2$CN$^{2+}$ collisions with neutrals.

Single electron transfer (SET), from the neutral to the dication, is a common outcome of dication-neutral collisions. Such dication-neutral electron transfer reactions are generally well represented by the Landau–Zener (LZ) ‘reaction window’ model.12 In this framework, an electron moves from the neutral to the dication at the point where the reactant (dication + neutral) and product (monocation + monocation) potential energy surfaces intersect. If this crossing lies within the ‘reaction window’, typically 2–6 Å, efficient electron-transfer can occur. At larger interspecies separations, there is weak coupling between the reactant and product potential energy surfaces and electron-transfer is disfavoured. For curve crossings at interspecies separations smaller than the ‘reaction window’, the coupling between the reactant and product potential energy surfaces is strong. In this situation it is likely that an electron will transfer twice between the neutral and dication, once on the neutral’s approach and once upon departure, resulting in no net electron transfer. Thus, when crossings lie within the reaction window, where the reactant and product potentials are not too strongly or too weakly coupled, efficient SET can occur. Typically, electron transfer exothermicities from 2–6 eV result in curve crossings that lie in the reaction window.92,93

Roithova et al.94 reported a comprehensive experimental study of the competition between PT and SET in the reactions of the hydrogen-containing dications CHX$^{2+}$ (X = F, Cl, Br, I) with several atomic, non-polar and polar neutral species including the neutral collision partners used in this study: Ar, N$_2$, and CO. From their observations, Roithova et al.94 proposed that if SET falls within the ‘reaction window’ (i.e. has an exothermicity of 2–6 eV) then PT is supressed, despite the latter pathway often nominally involving energetically favoured products. Roithova et al. argued that SET usually occurs at longer distances than PT, rendering the geometry of the collision less important, whereas PT requires at least a minimal degree of orientation of the reactants. Therefore, they deduced that ET is kinetically favoured. Roithova et al. also observed that when the neutral collision partner was polar, the relative intensity of PT increased. The authors accounted for this observation by invoking the increased attractive forces between the dication and polar neutral. These increased forces would result in an increased likelihood of a collision complex forming, thus extending the lifetime of the interaction and allowing more efficient access to the thermodynamic PT products. Despite the small dipole moment of CO, Roithova et al.94 surmised this dipole was reason for the increased relative intensity of PT with respect to ET in collisions with this molecule. Based on this study by Roithova et al.94 we might expect that, of the neutral collision partners used in our experiments (Ar, N$_2$, and CO), CO might exhibit the most intense PT channel.

As noted above, this work presents results from the experimental investigation of the bimolecular reactivity of CH$_2$CN$^{2+}$, a dication derived from the simplest nitrile, with the neutral species Ar, N$_2$, and CO. PT channels were observed in all the collision systems and these reactions are shown to proceed via direct, long-range, mechanism. The competition between the PT and SET reactions of CH$_2$CN$^{2+}$ does not follow the trends observed by Roithova et al.94 Specifically, SET does not occur following collisions between CH$_2$CN$^{2+}$ and Ar despite the presence of accessible pathways. Additionally, the relative intensity of the PT channel in the CH$_2$CN$^{2+}$ + CO system is lower than in the systems with non-dipolar collision partners. We also observe the presence of two CID channels, producing CH$_2$N$^+$ + C$^+$ and CHCN$^+$ + H$. In order to explain our results, we carried out electronic structure calculations of the CH$_2$CN$^{2+}$ dication, which reveal a complex energy landscape. Given the nature of the CID reactions observed, and the lowest energy structures revealed by our calculations, it appears that the CH$_2$CN$^{2+}$ dication involved in our experiments explores a number of conformations, including cyclic and linear isocyanide structures.

2. Experimental

Coincidence techniques involve the simultaneous detection of two or more products from a single reactive event. Dication
interactions with neutrals often generate pairs of monocations, and these pairs of ions are detected in coincidence in our position-sensitive coincidence (PSCO) mass spectrometry (MS) experiment. The apparatus used in this study has been described in detail in the literature. In brief, a pulsed beam of dications is directed into the field-free source region of a time-of-flight mass spectrometer (TOF-MS). In this region the dications interact with an effusive jet of the neutral reactant. Subsequent application of an extraction voltage to the source region allows the TOF-MS to detect the cation pairs generated from the dication-neutral interactions. The detection of these ions involves recording their arrival times, and arrival positions, at a large microchannel-plate detector. From this raw data, a list of flight times and arrival positions of the ions detected in pairs, a two-dimensional mass spectrum, can be generated, revealing the different reactive channels. The positional data accompanying the ionic arrivals yields the relative motion of the products of each detected event, providing a detailed insight into the mechanisms of each reactive channel.

The CH₂CN²⁺ ions used in the experiments described in this work were generated via electron ionization of CH₃CN (BDH, > 99.9%, purified by repeated freeze–pump–thaw cycles) by 100 eV electrons in a custom-built ion source. Previous studies have shown that whilst several molecular dications are generated following the bombardment of CH₃CN with electrons, CH₂CN²⁺ (m/z = 20) is the most abundant. Following electron ionization the positively charged ions are extracted from the ion source and pass through a hemispherical energy analyser to restrict the translational energy spread of the final CH₂CN⁺ beam to ~0.3 eV. The continuous beam of ions exiting the hemispherical analyser is then pulsed, using a set of electrostatic deflectors, before being accelerated and focussed into a commercial velocity filter. The velocity filter is set to transmit just the CH₂CN²⁺ (m/z = 20) ions. The resulting pulsed beam of energy-constrained CH₂CN²⁺ ions is then decelerated to less than 10 eV in the laboratory frame before entering the source region of the TOF-MS. In the TOF-MS source region the energy of dications is crossed with an effusive jet of the appropriate neutral species: Ar (BOC, 99.998%), N₂ (BOC, > 99.998%), or CO (Aldrich, > 99.0%). Single-collision conditions are achieved by employing an appropriately low pressure of the neutral collision partner. Under these conditions most dications do not undergo a collision and only a small percentage experience one collision. Such a pressure regime ensures no secondary reactions, due to successive collisions with two neutral species, influences the CH₂CN²⁺ reactivity we observe. As noted above, an electric field is applied across the TOF-MS source when the dication pulse reaches the centre of this region. This electric field accelerates positively charged species into the second electric field (acceleration region) of the TOF-MS and then on into the flight tube. At the end of the flight tube, the cations are detected by a position-sensitive detector comprising a chevron-pair of microchannel plates (diameter = 12.7 cm) located in front of a dual delay-line anode. The voltage pulse applied to the source region also starts the ion timing circuitry, to which the signals from the detector provide stop pulses. The experiments in this work employed a TOF-MS source field of 183 V cm⁻¹.

Signals from the detector are amplified and discriminated before being passed to a PC-based time-to-digital converter. If two ions are detected in the same TOF cycle, a coincidence event is recorded and each ion’s arrival time and impact position on the detector are stored for off-line analysis. The use of single-collision conditions ensures ‘false’ coincidences are kept to a minimum. The ion pairs data can be plotted as a 2D histogram, a ‘pairs spectrum’, where the time of flights (t₁, t₂) of each ion in the pair are used as the (x, y) co-ordinates. Peaks in the pairs spectrum readily identify bimolecular reaction channels that result in a pair of positively-charged product ions. Each such peak, the group of events corresponding to an individual reaction channel, can then be selected for further off-line analysis.

As shown in previous work, the positional and time of flight information for each ion of a pair can be used to generate their x, y and z velocity vectors in the laboratory frame; here the z-axis is defined by the principal axis of the TOF-MS. The x and y velocity vectors of an ion are determined from the associated positional information and the ion’s flight time; the z vector is determined from the deviation of the observed TOF from the expected TOF of the same ion with zero initial kinetic energy. The laboratory frame velocities are then converted into the centre-of-mass (CM) frame using the initial dication velocity. Often, the pair of monocations resulting from the reaction between a dication and a neutral are accompanied by a neutral species: a three-body reaction. A powerful feature of the PSCO-MS experiment is that the CM velocity of such a neutral product can be determined from the CM velocities of the detected ionic products via conservation of momentum.

To reveal the dynamics of a given reaction channel, a CM scattering diagram can be generated from the velocities of the product ions. Such CM scattering diagrams are radial histograms that, for each event collected in a given reaction channel, plot the magnitude of the products’ CM velocity |vᵢ| as the radial co-ordinate and the scattering angle θ between vᵢ and the CM velocity of the incident dication as the angular coordinate. In the kinematics that apply in our experiment, where the dication has significantly more momentum than the neutrals, the velocity of the incident dication is closely oriented with the velocity of the centre of mass. In our CM scattering diagrams, since θ ≤ θ ≤ 180°, the data for one product can be shown in the upper semi-circle of the figure and the data for another product in the lower semi-circle, as the scattering of each ion is azimuthally symmetric. The typical angular resolution of the scattering data achieved by the PSCO apparatus operating with a high source field is 4. The data presented in our two-dimensional figure on the page, we are integrating over the azimuthal angle: that is, we are binning events according to their value of θ irrespective of the azimuthal scattering angle. This integration over the azimuthal angle results in an isotropic scattering distribution yielding a sin(θ) distribution in our scattering diagrams, giving
our experimental arrangement a low relative detection efficiency at values of $\theta$ very close to 0° or 180°. This phenomenon is clearly visible in the scattering diagrams where even strongly forward scattered reactions exhibit a peak in $\theta$ away from $\theta = 0$. A more detailed discussion of the construction and form of these scattering diagrams can be found in the literature.\(^{96}\)

In all the scattering diagrams presented in this work, a logarithmic scale is used for the scattering intensity. Such scales allow the diagrams to reveal subtleties of the scattering in the areas with low intensities. However, one should be aware that such logarithmic scales can, at first glance, overemphasise the importance of areas in the scattering diagrams with low intensities. In the ESI† (Fig. S11) we contrast scattering diagrams with linear and logarithmic intensity scales.

For three-body reactions, internal-frame scattering diagrams can be a powerful aid in interpreting the reaction dynamics. In this class of scattering diagram $\tilde{\omega}$ is again the radial coordinate, but the angular coordinate is now the CM scattering angle with respect to CM velocity of one of the other product species.

From the CM velocities of the product species the total kinetic energy release (KER) $T$ for a given reactive event can also be determined using the individual CM velocities of the products.\(^{93}\) The exoergicity of the reaction $\Delta E$ can then be determined from $T$ and the CM collision energy, $E_{\text{com}}$:

$$\Delta E = T - E_{\text{com}} = -(E_{\text{products}} - E_{\text{reactants}}) \quad (1)$$

where $E_{\text{products}}$ and $E_{\text{reactants}}$ are the relative energies of the product and reactant states respectively. If the products lie lower in energy than the reactants, the resulting exoergicity will be positive. Performing this analysis for all the events collected for a given reaction channel provides a histogram of the exoergicities of the detected reactive events. From knowledge of the available electronic states of the reactants and products, the exoergicity spectrum can reveal the electronic states involved in the reaction.

### 2.1. Electronic structure calculations

Electronic structure calculations using Gaussian16 (Rev.A03)\(^{98}\) have been used to support the experimental findings. Specifically, the equilibrium geometries, and corresponding adiabatic and vertical ionisation energies, were determined for various species and structures of interest. These quantities were calculated for the lowest energy states of each of the relevant multiplicities. Stationary points were located using the MP2 algorithm with a cc-pVTZ basis set and the vibrational frequencies were analysed to confirm minima had been located. The energetics of these minima were then determined using single-point CCSD(T) calculations with the cc-pVTZ basis set. Zero-point energies from the MP2 geometry optimisations were used to correct the CCSD(T) single point energies. For the quartet isocyanate dication (4ID), where one MP2 vibrational frequency appeared unreliable, a B3LYP determination of the vibrational frequencies was used for the zero-point energy correction, the B3LYP algorithm converging to essentially the same geometry as the MP2 algorithm. The above methodology has been successfully used in previous work to determine the geometries and energetics of dications.\(^{88,99}\) Readers are referred to the SI for the full geometrical data for each dication structure.

### 3. Results and discussion

The collisions of CH$_2$CN$^{2+}$ with Ar, N$_2$ and CO were investigated, at CM collision energies of 5.0 eV, 4.5 eV, and 4.3 eV respectively, and the reactions observed are shown in Tables 1–3. In all the systems, a collision-induced dissociation (CID) channel was detected that involves the formation of CH$_2$N$^- + C^+$ (Rxns. I, IV, and X). Another CID channel, generating CHCN$^+ + H^+$, was also detected following collisions of CH$_2$CN$^{2+}$ with all three of the neutrals (Rxns. II, V, and XI). The collisions of CH$_2$CN$^{2+}$ with N$_2$ and CO also result in C$_2$N$^+ + H^+$ due to the further dissociation of the primary CID product HC$_2$N$^+$. (Rxns. VI and XII). All of the collision systems exhibit a PT channel forming HC$_2$N$^+$ and the protonated collision partner, XH$^+$ (X = Ar, N$_2$, CO, Rxns. III, VII, and XIII). Following collisions of CH$_2$CN$^{2+}$ with N$_2$ and CO, non-dissociative single electron-transfer (NDSET) channels are observed, forming CH$_2$CN$^+ + Y^+$ (Y = N$_2$, CO, Rxns. VIII and XIV). The collisions of CH$_2$CN$^{2+}$ with CO and N$_2$ also result in products from dissociative single electron-transfer (DSET) pathways, involving

| Table 1 | Reaction channels, following the collisions of CH$_2$CN$^{2+}$ with Ar at a CM collision energy of 5.0 eV, with associated relative intensities. See text for details |
| Reaction | Products | Relative intensity/% |
|---|---|---|
| I | CH$_3$N$^+ + C^+$ | 59.8 ± 3.3 |
| II | HC$_3$N$^+ + H^+$ | 22.0 ± 1.8 |
| III | HC$_3$N$^+ + ArH^+$ | 18.2 ± 1.6 |

| Table 2 | Reaction channels, following the collisions of CH$_2$CN$^{2+}$ with N$_2$ at a CM collision energy of 4.5 eV, with associated relative intensities. See text for details |
| Reaction | Products | Relative intensity/% |
|---|---|---|
| IV | CH$_3$N$^+ + C^+$ | 41.0 ± 0.6 |
| V | HC$_3$N$^+ + H^+$ | 15.6 ± 0.3 |
| VI | C$_2$N$^+ + H^+$ | 1.0 ± 0.1 |
| VII | HC$_3$N$^+ + N$_2$H^+ | 38.6 ± 0.5 |
| VIII | CH$_3$CN$^+$ + N$_2^+$ | 2.8 ± 0.1 |
| IX | HC$_3$N$^+ + N_2$ | 1.1 ± 0.1 |

| Table 3 | Reaction channels, following the collisions of CH$_2$CN$^{2+}$ with CO at a CM collision energy of 4.5 eV, with associated relative intensities. See text for details |
| Reaction | Products | Relative intensity/% |
|---|---|---|
| X | CH$_3$N$^+ + C^+$ | 21.3 ± 0.4 |
| XI | HC$_3$N$^+ + H^+$ | 3.7 ± 0.1 |
| XII | C$_2$N$^+ + H^+$ | 0.4 ± 0.1 |
| XIII | HC$_3$N$^+ + COH$ | 14.7 ± 0.3 |
| XIV | CH$_3$CN$^+ + CO^+$ | 52.4 ± 0.7 |
| XV | CO$^+ + CH^+$ | 2.3 ± 0.1 |
| XVI | HC$_3$N$^+ + CO^+$ | 0.8 ± 0.1 |
| XVIII | CH$_3$CN$^+ + C^+$ | 4.3 ± 0.2 |
the fragmentation of the CH$_2$CN$^+$ ion initially generated by SET (Rxns. IX, XV, and XVI). Finally, following the collisions of CH$_2$CN$^2+$ and CO, a DSET channel was detected, involving the fragmentation of CO$, generating CH$_2$CN$^+$ + C$^+$ (Rxn. XVII).

To determine the accessible electronic states of CH$_2$CN$^2+$ in the reactant beam the lowest energy structures of the CH$_2$CN$^2+$ dication were investigated using the computational methodology discussed above. Three local geometric minima of CH$_2$CN$^2+$ were discovered for both the doublet and quartet multiplicities (see Table 4). The lowest energy conformation is a doublet, with a cyclic C–C–N arrangement (2CD) shown in Fig. 1a, and an adiabatic double ionization energy of 28.9 eV from the ground state neutral. The ground state neutral structure determined from our calculations, a doublet with a linear geometry, agrees well with the literature. A linear doublet dication structure (2LD), shown in Fig. 1b, and a linear doublet isocyanide dication structure (2ID), shown in Fig. 1c, were also located. Each dication conformation also has an associated higher energy quartet state (Table 4). Whilst 2CD is the lowest energy conformation, a doublet with a linear geometry, agrees well with the literature. A linear doublet dication structure (2LD), shown in Fig. 1b, and a linear doublet isocyanide dication structure (2ID), shown in Fig. 1c, were also located. Each dication conformation also has an associated higher energy quartet state (Table 4). Whilst 2CD is the lowest energy conformation revealed by our calculations, the linear and linear isocyanide structures do not lie far above the 2CD state energetically. Therefore, it is quite possible that the CH$_2$CN$^2+$ ions in our beam are fluxional, sampling different minima on the dicaticonic potential energy surface. Here we can directly analogue with the benzene dication which has a number of metastable minima on the dicaticonic potential energy surface which are sampled by C$_6$H$_6$$^2+$ ions.

The six metastable dication states revealed by the calculations (Table 4) have adiabatic double ionization potentials, from CH$_2$CN, of between 28.9 eV and 32.9 eV. These double ionization potentials are significantly lower than the double ionization potentials of Ar (43.39 eV), $^1$N$_2$ (~ 43 eV), and CO (~ 41.5 eV). Therefore, as expected, there are no product peaks in any of the spectra from the CH$_2$CN$^2+$ collision systems that result from double electron-transfer (DET), as such processes are significantly endoergic.

Similar computational investigations were also undertaken to determine likely structures of the CH$_2$CN$^+$ monocation, the ion which results from the dication acquiring an electron in a SET reaction. These calculations reveal a singlet and triplet state for each of the cyclic, linear, and linear isocyanide geometries, shown in Table 5. The vertical ionization energies determined for the linear and linear isocyanide geometries are a good match (~ 0.1 eV difference) with experimental data from photoelectron spectroscopy. In addition the energetics and geometries we derive are in good agreement with recent computational studies.

Calculations were also employed to determine the energy of the lowest energy structure of the HC$_2$N$^+$ monocation formed when the dication loses a proton. The lowest energy state of HC$_2$N$^+$ determined from these calculations is a doublet that has a linear N–C–C geometry, with the H bonded to the terminal C slightly off the N–C–C axis. This state has a calculated adiabatic ionization energy of 10.4 eV. The structure of the corresponding neutral species, HCCN, was also determined to corroborate our methodology. This investigation revealed a staggered configuration, in accord with the geometry calculated by Nimlos et al.

In the following sections the different reaction mechanisms revealed by our scattering data for the collisions of CH$_2$CN$^2+$ + X (X = Ar, N$_2$, and CO) are discussed.

### 3.1. The formation of C$^+$ + CH$_2$N$^+$ (CID)

Fig. 2 shows the CM scattering diagrams of the C$^+$ and CH$_2$N$^+$ products resulting from the collisions of CH$_2$CN$^2+$ with Ar, N$_2$, and CO. It is clear, from the dependence of the number of

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**Table 4** Abbreviations and adiabatic ionization energies of the metastable CH$_2$CN$^2+$ dication structures located by the computational investigation. See text for details.

| Multiplicity | Cyclic conformation | Linear conformation | Linear isocyanide conformation |
|--------------|---------------------|---------------------|-------------------------------|
| Doublet      | 2CD, 28.9 eV        | 2LD, 30.4 eV        | 2ID, 29.4 eV                  |
| Quartet      | 4CD, 32.9 eV        | 4LD, 32.7 eV        | 4ID, 32.1 eV                  |

**Table 5** Abbreviations and adiabatic ionization energies of the metastable CH$_2$CN$^+$ monocation structures, see text for details.

| Multiplicity | Cyclic conformation | Linear conformation | Linear isocyanide conformation |
|--------------|---------------------|---------------------|-------------------------------|
| Singlet      | 1CM, 11.0 eV        | 1LM, 10.1 eV        | 1IM, 10.2 eV                  |
| Triplet      | 3CM, 12.9 eV        | 3LM, 12.1 eV        | 3IM, 12.3 eV                  |

**Fig. 1** Geometries of the local minima of the reactant CH$_2$CN$^2+$ dication (doublet states). (a) Cyclic conformation, 2CD, (b) linear conformation, 2LD, (c) linear isocyanide conformation, 2ID. See text for details.
counts in this channel on the neutral target gas pressure, that these products result from bimolecular reactions. The scattering diagrams (Fig. 2) for each of the reactions are remarkably similar, with the C’ and CH₂N⁺ ionic products isotropically scattered about the CM velocity of the reactant dication. Therefore, the mechanism responsible for this channel is clearly CID: the CH₂CN₂⁺ dication dissociates as a result of a collision with the neutral species, forming CH₂N⁺ + C’. CID involves energy transfer between the dication and neutral species promoting the dication to a state which subsequently dissociates. We are unable to directly resolve the energy transfer between the dication and neutral species in this CID channel because we cannot directly determine the change in kinetic energy of the neutral product. However, determination of the CM velocity associated with the ionic products reveals, within our experimental uncertainty, no resolvable difference to the velocity of the reactant dication. This observation clearly indicates there is no energy transfer between the dication and neutral species. Interestingly, modelling shows that CH₂N⁺, the product of this CID reaction, is one of the most abundant ion species in the atmosphere of Titan.117–119

The precursor used to produce CH₂CN₂⁺ via electron ionization is acetonitrile, CH₃CN. CH₃CN has a well-defined structure (C₆H₅) involving a –CH₃ moiety bound to the nitrile group (–CN). Given this connectivity, the presence of a CID channel generating C’ (CH₂N⁺ + C’) points to the involvement of a dication structure significantly different from the structure of the CH₂CN precursor, which has no terminal carbon atoms. As discussed above, the minimum energy CH₂CN₂⁺ structure, 2CD, revealed by the calculations has a cyclic C–C–N arrangement (Fig. 1a), whilst slightly higher in energy is the linear isocyanide structure, 2ID (Fig. 1c). It is easy to see how either the 2CD or 2ID CH₂CN₂⁺ structures could readily dissociate to form CH₂N⁺ + C’. The presence of this CID channel therefore hints that our dication beam includes dications in the cyclic or linear isocyanide geometries. Of course, an alternative explanation could be that upon collision with the neutral, the linear dication rearranges to the cyclic or isocyanide state before dissociation. Our data cannot reveal which of these detailed mechanisms is operating. However, the similarity in scattering across the three CH₂CN₂⁺ + neutral systems strongly suggests the cyclic or linear isocyanide dication structures are involved. Extracting the ionic velocity distributions from our data for this channel, we see that velocity distributions of the CH₂N⁺ and C’ products are independent of the neutral species involved in the collision, as shown qualitatively in Fig. 2. As well as the dominant isotropic angular distribution, the scattering diagrams (Fig. 2) have a slight ‘bump’ at angles near 90 degrees, where the product ions are scattered with higher velocities. This bump is also reflected in the experimental exoergicity spectrum (Fig. 3, discussed below), giving rise to the structure from ~8–11 eV.

Fig. 3 shows the experimental exoergicity distributions for the dissociation of CH₂CN₂⁺ in these CID channels forming CH₂N⁺ + C’. The exoergicity distributions for the different collision systems are again very similar, rising sharply from ~5 eV to a maximum at ~6 eV before slowly falling. The experimental exoergicity distributions all have FWHMs of ~5.5–7.5 eV, with a shoulder to higher exoergicities. In order to rationalise the experimental exoergicity distribution we must consider the relative energies of the reactant and product species that could be involved in this reaction. For this system, energetic data for C’ is readily available and the relevant energy levels for CH₂N⁺ can be estimated using the proton affinity and heat of formation of HCN, which are also in the literature.120,121

By combining the energy of a proton, the heat of formation of HCN and the proton affinity of HCN, we determine an estimate of the heat of formation of CH₂N⁺ as 9.87 eV and will assume that no other electronic states of CH₂N⁺ are involved in the reaction. For simplicity, and given our experimental observations above, we will assume that, given the level of accuracy of these energetic estimates, there is negligible energy transfer between the dication and neutral species. With the above energetics, and using the double ionization energies from our calculations of CH₂CN₂⁺ and the heat of formation of CH₂CN (there is some uncertainty in this latter value, we will use +2.5 eV),122–124 we are then in a position to estimate the literature exoergicities expected for this reaction. These energetics indicate that, in an exothermic reaction, C’ could be formed in either its ground, ³P, or first excited state, ⁴P.

Using the above energetics, we find three pathways with literature exoergicities that fall within the FWHMs observed experimentally for these CID channels. The energy releases from 4CD and 4ID dissociating to produce C’ + CH₂N⁺ in their electronic ground states are 6.9 eV and 6.0 eV respectively,
fitting nicely with the experimental exoergicity distributions in Fig. 3. The corresponding energy release from the linear quartet CH$_2$CN$^{2+}$ dication is 6.7 eV, also falling within the observed experimental exoergicity distribution. However, due to the fragments formed we think that the dication reactant is likely not in the regular linear conformation (4LD). If either of the doublet 2CD (the ground state) or 2ID states were involved, the expected exoergicity released from this CID channel would be 2.8 eV and 3.3 eV respectively, clearly outside of the observed experimental exoergicity range (Fig. 3). Of course, both the reactant CH$_2$CN$^{2+}$ and product CH$_2$N$^+$ ions could be vibrationally excited, which would act to broaden the experimentally observed exoergicities. If the reactant dication was in the 2CD state and significantly vibrationally excited (on the order of ~3 eV) the resulting exoergicities would be a match with the observed distribution. However, this level of vibrational excitation is very unlikely as such dications would almost certainly dissociate before reaching the interaction region. Further detailed computational work, beyond the scope of this study, including their detailed computational work, beyond the scope of this study, would be valuable to further interpret the exoergicity data.

To summarise, one of the dominant channels following the collisions of CH$_2$CN$^{2+}$ with Ar, N$_2$, and CO results in the generation of CH$_2$N$^+$ + C$^+$, revealed by the dynamics to be due to CID. This CID channel is the most intense reaction channel following the collisions of CH$_2$CN$^{2+}$ with Ar and N$_2$, whilst with CO, it is the second most intense channel. The presence of this CID channel, resulting in the ejection of a C$^+$ from the reactant dication, suggests that the dications participating in this CID channel are likely not in the linear H$_2$CCN conformation derived directly from the structure of the CH$_2$CN precursor. The experimental exoergicity distributions of these three channels are in excellent agreement with each other and point to the involvement of either the quartet cyclic state, 4CD, or the quartet linear isocyanide state, 4ID, of the CH$_2$CN$^{2+}$ dication, and result in the formation of C$^+$ + CH$_2$N$^+$ in their electronic ground states. The CH$_2$N$^+$ product is likely formed with some vibrational excitation.

3.2. The formation of H$^+$ + HC$_2$N$^+$ (CID)

Each of the coincidence spectra resulting from the CH$_2$CN$^{2+}$ + X (X = Ar, N$_2$, CO) collision systems displays a peak corresponding to the formation of a pair of ions with $m/z = 1$ and $m/z = 39$, corresponding to H$^+$ + HC$_2$N$^+$, clearly another CID channel. The form of the pairs spectra and the associated scattering diagrams from the channels generating HC$_2$N$^+$ + H$^+$ in each of the three CH$_2$CN$^{2+}$ + X (X = Ar, N$_2$, CO) systems look close to identical, pointing to the involvement of the same mechanism in each system (see Fig. SI2 and SI3 in the ESI†). It is difficult to say too much about the mechanism of this channel from the dynamics because the velocity of the H$^+$ product ion is dominant, leading to markedly increased uncertainty in the scattering of the molecular ion; however, the form of the scattering (with product monocation distributions centred on the dication velocity) clearly indicates that CID is the mechanism involved.

The experimental exoergicity distributions recorded following the CH$_2$CN$^{2+}$ + X → HC$_2$N$^+$ + H$^+$ + X (X = Ar, N$_2$, CO) reactions are shown in Fig. 4. All three exoergicity distributions have a similar structure, with a peak centred at ~2.5 eV and a FWHM from approximately 0.5 eV to 6 eV. In order to rationalise these experimental exoergicities, we must again consider the product and reactant states involved. As before, we will consider the six metastable states of CH$_2$CN$^{2+}$ found from our calculations: 2CD, 2LD, 2ID, 4CD, 4LD, and 4ID. The neutral reactants will, of course, be in their ground vibronic states. As before, given our experimental observations, we will neglect any energy transferred to the dication in the collision. The heat of formation of H is +2.26 eV, and its ionization energy is 13.60 eV. The lowest energy structure of the HC$_2$N$^+$ monocation revealed by our calculations (a doublet with a linear N–C–C geometry) has an adiabatic ionization energy of 10.4 eV relative to the neutral, and the neutral has a heat of formation of +4.8 eV. The literature exoergicities of the reaction CH$_2$CN$^{2+}$ + X → HC$_2$N$^+$ + H$^+$ + X (X = Ar, N$_2$, CO) using the above constraints are marked on the experimental exoergicity distributions in Fig. 4 for each of the six possible CH$_2$CN$^{2+}$ dication states. These exoergicities are 2CD = 0.4 eV,
result from vibrational excitation of the HCCN+ product. It could involve reactions of the 2CD and 2ID geometries or could occur in this CID channel. The signals at lower exoergicities in Fig. 4 are an example of, predominantly, the linear 2LD state and the three quartet states (4CD, 4ID, 4LD) are good matches to the bulk of the experimental data (Fig. 4). It is tempting to attribute the linear dication structures 2LD and 4LD to this CID reaction, given that we think that these structures are not involved in the CID channel forming C+ + CH2N+. However, the data does not rule out the involvement of the 4ID and 4CD dication states in this CID channel. The signals at lower exoergicities in Fig. 4 could involve reactions of the 2CD and 2ID geometries or could result from vibrational excitation of the HCCN+ product.

To summarise, the peak observed in the coincidence spectrum corresponding to the formation of HC2N+ + H+ following the collisions of CH2CN2+ + N2/CO reveal a clear additional peak corresponding to the formation of C2N+ (m/z = 38) + H+ (m/z = 1). This channel likely also results in the formation of the neutrals N2/CO + H. The velocities of the ionic products are similar to those for the above CID reactions that result in the formation of CHCN+ + H+, indicating a similar mechanism is responsible. However, since this reaction results in the formation of more than one neutral species, it is hard to probe this reaction in more detail using our experimental methodology. However, given the above observations, it seems clear that these channels also proceed via CID forming H+ + HC2N+ followed by the dissociation of the nascent HC2N+ to form C2N+ + H. Indeed, the potential vibrational excitation of the primary HC2N+ product was remarked upon on above.

3.2.1. The formation of H+ + C2N+. The coincidence spectra recorded following the collisions of CH2CN2+ + N2/CO reveal a clear additional peak corresponding to the formation of C2N+ (m/z = 38) + H+ (m/z = 1). This channel likely also results in the formation of the neutrals N2/CO + H. The velocities of the ionic products are similar to those for the above CID reactions that result in the formation of CHCN+ + H+, indicating a similar mechanism is responsible. However, since this reaction results in the formation of more than one neutral species, it is hard to probe this reaction in more detail using our experimental methodology. However, given the above observations, it seems clear that these channels also proceed via CID forming H+ + HC2N+ followed by the dissociation of the nascent HC2N+ to form C2N+ + H. Indeed, the potential vibrational excitation of the primary HC2N+ product was remarked upon on above.

3.3. Proton-transfer (PT) channels

PT is observed in all three CH2CN2+ + neutral collision systems, resulting in the formation of HC2N+ and the protonated neutral. PT is often observed following the collisions of hydrogen-containing dications with neutrals, and has been detected before in different reactive systems involving each of the neutral species used in this investigation: N2, CO, and Ar. Considering the products of PT, the argonium ion, ArH+, is an example of a molecular species involving a rare gas and has been detected in the ISM, including in other galaxies. N2H+ has been detected in the interstellar medium, protoplanetary disks, and in the atmosphere of Titan (where modelling also predicts the presence of N2H+). In such environments N2H+ can be formed via a number of processes. OCH− has also been detected in the ISM.

Fig. 5 shows the CM scattering diagrams of the ionic products resulting from the PT reactions. In each of the scattering diagrams, strong forward scattering is observed. Such a scattering pattern is indicative of a long-range, direct, mechanism. If a proton transfers from the dication to the neutral at a large interspecies separation, the deflection of the reactant species is minimal and therefore the product velocities are strongly oriented with the velocities of their corresponding reactants. Evidence shows that dication-neutral reactions involving the formation of new bonds more commonly occur via the formation of a collision complex, where the reactants temporarily associate in order for new bonds to be made. However, previous experiments have also shown that dication-neutral PT reactions can occur via direct mechanisms. For example, PT occurs via a direct mechanism following collisions of C2H3+ with Ar. The direct nature of these PT reactions is not surprising given that just a proton is being transferred between the reactants, and thus the operation of a direct mechanism, analogous to LZ style electron transfer, seems dynamically reasonable.

The experimental exoergicity distributions recorded for the PT channels are shown in Fig. 6. The distributions have similar...
shapes for all three collision systems, with maxima at ~5 eV. In order to rationalise the experimental exoergicity spectra we again must consider the reactant and product electronic states that could be involved. We will again consider the six CH2CN2+ dication states obtained from our calculations (Table 4), and use the adiabatic ionization energy from our calculations of lowest energy structure of the HC2N+ monocation, together with the heat of formation from Nimlos et al.116 The energetics of the protonated species ArH+ and N3H+ are easily determined from proton affinities that are readily available, giving heats of formation of 12.0 eV and 10.7 eV respectively.104 CO can be protonated at either the carbon, forming OCH+, or the oxygen, forming COH+, giving heats of formation of 8.6 eV and 10.3 eV respectively.104

Using these energetics, the literature exoergicities for the PT reactions involving the cyclic doublet dication 2CD, resulting in the formation of the XH+ and HC2N+ products in their ground states are 4.2 eV, 5.5 eV, 4.8 eV and 6.6 eV for the reactions with Ar, N2, CO (forming COH+) and CO (forming OCH+) respectively. These literature exoergicities are marked on the corresponding distributions in Fig. 6 and fit nicely with the experimental data. Of course, if the linear isocyanide dication structure was involved (2ID), similar exoergicities (~0.4 eV higher) would result. Therefore, for simplicity, only 2CD is discussed below. Predicted literature exoergicities involving the quartet dication states are ~3 eV higher than the literature exoergicities shown in Fig. 5, clearly outside of the range of the experimental exoergicity distributions. Therefore, the involvement of the quartet states in this PT channel is, at most, minor. As seen in Fig. 6c, the literature exoergicity predicted for the formation of COH+ is a better fit with the experimental exoergicity distribution than the formation of the lower energy OCH+ isomer.

If the reactant dication was vibrationally excited, this would act to increase the observed experimental exoergicity. Conversely, if the product HC2N+ or XH+ ions were formed with vibrational excitation, the observed experimental exoergicity would be lower. So given the likely broadening of the experimental distributions due to such vibrational excitation, the experimental data (Fig. 6) can clearly be accounted for with the involvement of the 2CD (and 2ID) dication geometries. The experimental exoergicity distributions (Fig. 6) are clearly broader for the molecular targets, N2 and CO. Such broadening, particularly to low exoergicities is likely due to the increased number of vibrational modes in the triatomic (N3H+, COH+) product ion facilitating vibrational excitation of the XH+ product in comparison with the diatomic ArH+.

To summarise, PT reactions occur in each of the CH2CN2+ + X (X = Ar, N2 and CO) systems, resulting in the formation of HC2N+ and XH+. The dynamics show that PT occurs via a direct, long-range, mechanism. The experimental exoergicity distributions reveal that the 2CD and 2ID dication states are the
principal reactants, generating the monocationic products in their ground electronic states.

3.4. Non-dissociative single electron-transfer (NDSET) channels

Fig. 7 shows the CM scattering diagrams for the NDSET reactions \( \text{CH}_2\text{CN}^{2+} + X \rightarrow \text{CH}_2\text{CN}^+ + X^+ \), for \( X = \text{N}_2 \) and CO. Strong forward scattering is observed for both reactions. That is, the velocity of the \( \text{CH}_2\text{CN}^- \) product ion is oriented with the velocity of the reactant dication, \( \text{CH}_2\text{CN}^{2+} \). Such a scattering pattern is typical for NDSET reactions in dication-neutral systems and is indicative of a direct LZ electron-transfer process occurring at a significant interspecies separations.\(^{65,72,92,97,138}\) For the NDSET reaction with \( \text{N}_2 \), in the coincidence spectrum only the forward scattered counts could be isolated from counts from the PT channel, hence the absence of signal at higher angles in the scattering diagram (Fig. 7a).

Fig. 8 and 9 show the experimental exoergicity distributions recorded following the NDSET reactions of \( \text{CH}_2\text{CN}^{2+} \) with \( \text{N}_2 \) and CO respectively. As above, in order to rationalise the experimental exoergicity spectra we must first consider the reactant and product states that could be involved. Again, we consider the six electronic states of \( \text{CH}_2\text{CN}^{2+} \) (Table 4). Calculations probing the structure of \( \text{CH}_2\text{CN}^+ \) reveal a singlet and triplet state for each of cyclic, linear, and linear isocyanide geometries (See Table 5): 1CM, 1LM, 1IM, 3CM, 3LM, and 3IM. The lowest energy \( \text{CH}_2\text{CN}^- \) state is the singlet linear conformation (1LM). The ground state of \( \text{N}_2^- (X^2\Sigma_g^+) \) is 15.58 eV higher in energy than the ground state of \( \text{N}_2 \).\(^{106,139}\) The lowest energy dissociation asymptote of \( \text{N}_2 (\text{C}_2\Sigma_u^+ + \text{N}(\text{4S})) \) lies at \( \sim 24.3 \) eV above ground state \( \text{N}_2 \). The energy of this dissociation asymptote corresponds to the energy of \( \text{N}_2^+ (\text{C}_2\Sigma_u^+ + \nu = 3) \).\(^{104,140}\) Photoionization studies of \( \text{N}_2 \) show that \( \text{N}_2^+ \) states generated at an energy higher than 24.3 eV have dissociation lifetimes significantly less than the timescale of our experiment and therefore will not contribute to the observed \( \text{N}_2^+ \) counts in this channel.\(^{140–143}\) The ground state of \( \text{CO}^+ (X^2\Sigma_g^-) \) lies 14.0 eV above \( \text{CO}(X^2\Sigma_g^-) \).\(^{144,145}\) The lowest energy dissociation asymptote of \( \text{CO}^+ \), corresponding to \( \text{C}(\text{2P}) + \text{O}(\text{3P}) \), lies at 22.4 eV above \( \text{CO}(X^2\Sigma_g^-) \).\(^{104}\) Photoionization investigations of \( \text{CO} \) also show that if \( \text{CO}^+ \) is formed at an energy of greater than 22.4 eV, it will dissociate within the lifetime of the PSCO experiment and therefore will not contribute to the counts observed in this reaction channel.\(^{145–147}\)

To use the above energetic information to analyse the experimental exoergicity distributions and determine the electronic states involved in the NDSET channels we will also assume that the transitions will be spin allowed, and that the geometry of the dication will not change upon accepting an electron; for example, a cyclic dication must result in the formation of a cyclic monocation when it accepts an electron. Indeed, experimental work shows “vertical” transitions dominate dicationic electron transfer.

The experimental exoergicity distribution for the NDSET reaction with \( \text{N}_2 \) (Fig. 8) has a maximum at 4.0 eV, with a FWHM from 3.6–5.2 eV. Using the energetics and assumptions presented above, we find that there are 9 pathways that have literature exoergicities that match the experimental data, which we list in Table 6. The exoergicities of these channels are also...
within the LZ reaction window. These available pathways reaction with Ar is surprising as there are several possible are no experimental signals associated with SET reactions in our calculations. Hence there are no energetic grounds for the absence of NDSET reactivity following collisions of CH$_2$CN with Ar. The absence of any SET channels in the CH$_2$CN$^{+3}$ + Ar collision system, despite energetically accessible product channels, clearly points to a kinetic barrier in this pathway. Indeed, ‘Coulomb barriers’ on dication/neutral interaction potentials have been proposed to account for the stability of both molecular dication complexes, preventing charge separation by the routes of CID and SET. In a dication-neutral collision system such a Coulomb barrier could impede access to the SET product asymptote. For the CH$_2$CN$^{+3}$ + Ar collision system, we see that SET is supressed, but PT is efficient. Thus, any barrier is clearly only significant in the SET exit channel. There is further experimental evidence supporting the existence of a kinetic barrier to the SET products. Specifically, we do observe signals, in the coincidence spectra, for SET between CH$_2$CN$^{+3}$ with Ar in collisions beyond the interaction region of our TOF-MS; when the dications are moving much faster and the collision energy is correspondingly higher. These higher collision energies clearly allow the barrier to SET in the collision system with Ar to be overcome. Indeed, despite the observation by Roithová et al. that the involvement of polar neutral collision partners increases the chance of PT reactions occurring, in the current study the CH$_2$CN$^{+3}$ + CO system exhibits the highest ratio of SET to PT reactivity. Clearly there are many subtleties in individual dication-neutral interaction potentials that can undermine global generalizations regarding dicationic reactivity.

3.4.1. Dissociative single electron-transfer (DSET) channels with CO and N$_2$. Examination of the dynamics of the product ion velocities of the remaining peaks in the coincidence spectra recorded following the collisions of CH$_2$CN$^{+3}$ with N$_2$ and CO, reveal that there are four channels resulting from DSET. Specifically, the pairs spectrum shows peaks for N$_2$/$\text{CO}^+$ (m/z = 28) + HC$_2$N$^+$ (m/z = 39), where the nascent CH$_2$CN$^{+3}$ ion formed following initial SET loses an H atom. There is also a peak for CO$^+$ (m/z = 28) + CH$^+$ (m/z = 13), resulting from the fragmentation of the CH$_2$CN$^+$ primary SET product. Additionally, there is a channel involving the formation of C$^+$ (m/z = 12) + CH$_2$CN$^+$ (m/z = 40), where the CO$^+$ generated by initial SET then fragments into C$^+$ + O. These DSET channels have low intensities. For example, the DSET channel in the CH$_2$CN$^{+3}$ + N$_2$ system contributes just 1.1% of the total counts. In order to dissociate, accessing the C$^+$ + O asymptote, CO$^+$ must be formed with an energy of at least 22.4 eV relative to CO. Using the energies of the CH$_2$CN$^{+3}$ dication states and CH$_2$CN$^+$ monocation states we determined computationally, we see that forming CO$^+$ at such an energy is possible but requires the formation of a singlet CH$_2$CN$^+$ state from a quartet CH$_2$CN$^{+3}$ state, which is nominally spin forbidden if forming a doublet state of CO$^+$. This spin forbidden nature likely accounts for the very low intensity of this DSET channel. In order for CO$^+$ to fragment to produce O$^+$ + C, which we do not observe experimentally, CO$^+$ must be formed with an energy at least 24.6 eV above that of neutral CO. Our calculated energetics show that it

| Pathway | Reactant CH$_2$CN$^{+3}$ state | Product CH$_2$CN$^+$ state populated | Product CO$^+$ state populated | Literature exoergicity (eV) |
|---------|-------------------------------|-------------------------------------|-------------------------------|-----------------------------|
| (j)     | 2CD                          | 1CM                                 | X$\Sigma^+$                  | 3.8                         |
| (k)     | 4CD                          | 3CM                                 | X$\Sigma^+$                  | 6.1                         |
| (l)     | 2LD                          | 1LM                                 | X$\Sigma^+$                  | 3.5                         |
| (m)     | 2LD                          | 1LM                                 | X$\Sigma^+$                  | 6.3                         |
| (n)     | 2LD                          | 1LM                                 | A$\Pi$                       | 3.7                         |
| (o)     | 2LD                          | 3LM                                 | X$\Sigma^+$                  | 4.3                         |
| (p)     | 4LD                          | 3LM                                 | X$\Sigma^+$                  | 6.6                         |
| (q)     | 4LD                          | 3LM                                 | A$\Pi$                       | 4.0                         |
| (r)     | 2LD                          | 1LM                                 | X$\Sigma^+$                  | 5.1                         |
| (s)     | 2LD                          | 3LM                                 | X$\Sigma^+$                  | 3.1                         |
| (t)     | 4LD                          | 3LM                                 | X$\Sigma^+$                  | 5.8                         |
| (u)     | 4LD                          | 3LM                                 | A$\Pi$                       | 3.2                         |

Table 6: Electronic states involved, with associated literature exoergicities, for the reaction CH$_2$CN$^{+3}$ + N$_2$ + $\text{O}^{12}\text{C}$ → CH$_2$CN$^+$ + N$_2$ + $\text{O}^{12}\text{C}$.
is not possible to form CO at such an energy from the \( \text{CH}_2\text{CN}^{2+} \) and \( \text{CH}_2\text{CN}^- \) states we believe to be involved in this collision system. Such agreement between the conclusions from our energetics and the experimental observations strongly supports the robustness of our computational conclusions concerning the accessible \( \text{CH}_2\text{CN}^{2+} \) and \( \text{CH}_2\text{CN}^- \) states.

4. Conclusions

In this work we have explored the reactivity of the molecular dication \( \text{CH}_2\text{CN}^{2+} \) with \( \text{Ar}, \text{N}_2, \) and \( \text{CO} \). To the author’s knowledge, this is the first study of the bimolecular reactivity of \( \text{CH}_2\text{CN}^{2+} \), a species relevant to the ionospheres of planets and satellites including Titan. Calculations of the lowest energy electronic states of \( \text{CH}_2\text{CN}^{2+} \) reveal six low-lying (and presumably metastable) bound electronic states. Specifically, we locate both doublet and quartet states for cyclic (2CD & 4CD), linear (2LD & 4LD), and linear isocyanide (2ID & 4ID) geometries. The lowest energy electronic state was determined to be the doublet cyclic dication, 2CD with an adiabatic double ionization energy of 28.9 eV.

All of the collision systems we investigated display an intense collision induced dissociation (CID) channel resulting in the formation of \( \text{C}^+ + \text{CH}_2\text{N}^+ \). This channel exhibits similar dynamics and experimental exoergicity distributions for all three collision systems. The formation of a \( \text{C}^+ \) fragment suggests the involvement of the cyclic or linear isocyanide dication states, and the experimental exoergicity distributions confirm the involvement of the quartet cyclic or isocyanide (4CD, 4ID) dication states.

Each \( \text{CH}_2\text{CN}^{2+} \) + neutral collision system also results in a CID reaction generating \( \text{H}^+ + \text{HC}_2\text{N}^- \). Again, the products exhibit similar dynamics and experimental exoergicity distributions in all three collision systems. The experimental exoergicity distributions primarily point to the involvement of the 2CD and 2ID dication states in this CID reaction, with 2LD also involved. Channels involving further dissociation of the \( \text{HC}_2\text{N}^- \) product, forming \( \text{H}^+ + \text{CCN}^- \), were also detected in the \( \text{CH}_2\text{CN}^{2+} \) + N\(_2\)/CO systems.

Proton transfer (PT) occurs in all three of the collision systems via a direct mechanism, giving further evidence that hydrogen-containing dications are effective proton donors. The experimental exoergicity distributions resulting from the PT reactions firmly point to the involvement of the 2CD or 2ID dication states as the reactants.

Finally, we observe reaction channels involving single-electron transfer (SET) following the collisions of \( \text{CH}_2\text{CN}^{2+} \) with both \( \text{N}_2 \) and \( \text{CO} \). The experimental exoergicity distributions show that the non-dissociative SET channels could involve any of the dication states in the beam. There is no SET reactivity with \( \text{Ar} \) which we attribute, with experimental evidence, to an energy barrier restricting access to the appropriate product asymptote. In this work, SET channels are only more intense than PT processes in collisions with the CO molecule. These observations differ from the trends observed by Roithová et al.\(^{24}\) for reactions of halogenated dications.

Our experimental data indicates the reactivity we observe involves all three dication conformations (cyclic, linear, isocyanide) identified by our computational investigation of the accessible structures of \( \text{CH}_2\text{CN}^{2+} \). This observation points towards considerable fluxionality in the structure of the low-lying metastable electronic states of \( \text{CH}_2\text{CN}^{2+} \).

Conflicts of interest

There are no conflicts to declare.

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References

1. H. S. Bridge, J. W. Belcher, A. J. Lazarus, J. D. Sullivan, R. L. McNutt, F. Bagenal, J. D. Scudder, E. C. Sittler, G. L. Siscoe, V. M. Vasyliunas, C. K. Goertz and C. M. Yeates, Science, 1979, 204, 987–991.
2. S. Ghosh, K. K. Mahajan, J. M. Grebowsky and N. Nath, J. Geophys. Res., 1995, 100, 23983–23991.
3. J. H. Hoffman, C. Y. Johnson, J. C. Holmes and J. M. Young, J. Geophys. Res., 1969, 74, 6281–6290.
4. E. Dubinin, R. Modolo, M. Fraenz, J. Woch, G. Chanteur, F. Duru, F. Akalin, D. Gurnett, R. Lundin, S. Barabash, J. D. Winningham, R. Frahm, J. J. Plaut and G. Picardi, J. Geophys. Res.: Space Phys., 2008, 113, A10217.
5. J. Lilensten, C. Simon, O. Witasse, O. Dutuit, R. Thissen and C. Alcaraz, Icarus, 2005, 174, 285–288.
6. L. A. Frank, W. R. Paterson, K. L. Ackerson, V. M. Vasyliunas, F. V. Coroniti and S. J. Bolton, Science, 1996, 274, 394–395.
7. H. Gu, J. Cui, D. Niu, L. Dai, J. Huang, X. Wu, Y. Hao and Y. Wei, Earth Planet. Phys., 2020, 4, 396–402.
8. C. Simon, J. Lilensten, O. Dutuit, R. Thissen, O. Witasse, C. Alcaraz and H. Soldi-Lose, Ann. Geophys., 2005, 23, 781–797.
9. O. Witasse, O. Dutuit, J. Lilensten, R. Thissen, J. Zabka, C. Alcaraz, P.-L. Bletty, S. W. Bougher, S. Engel, L. H. Andersen and K. Seiersen, Geophys. Res. Lett., 2002, 29, 104.
10. J. Lilensten, O. Witasse, C. Simon, H. Soldi-Lose, O. Dutuit, R. Thissen and C. Alcaraz, Geophys. Res. Lett., 2005, 32, L03203.
11. A. Beth, K. Altwegg, H. Balsiger, J.-J. Berthelier, M. R. Combi, J. De Keyser, B. Fiethé, S. A. Fuselier, M. Galand, T. I. Gombosi, M. Rubin and T. Sémon, Astron. Astrophys., 2020, 642, A27.
12. S. D. Price, J. D. Fletcher, F. E. Gossan and M. A. Parkes, Int. Rev. Phys. Chem., 2017, 36, 145–183.
145 Y. Zhao, M. Cao, Y. Li, X. Shan, F. Liu, L. Sheng, L. Li and W. Liu, J. Electron Spectrosc. Relat. Phenom., 2014, 196, 181–186.

146 P. Baltzer, M. Lundqvist, B. Wannberg, L. Karlsson, M. Larsson, M. A. Hayes, J. B. West, M. R. F. Siggel, A. C. Parr and J. L. Dehmer, J. Phys. B: At., Mol. Opt. Phys., 1994, 27, 4915–4932.

147 R. C. Shiell, M. Evans, S. Stimson, C.-W. Hsu, C. Y. Ng and J. W. Hepburn, Chem. Phys. Lett., 1999, 315, 390–396.

148 J. Roithová and D. Schröder, Angew. Chem., Int. Ed., 2009, 48, 8788–8790.

149 J. F. Lockyear, PhD thesis, UCL, 2011.