Evidence of polarisation in the prompt gamma-ray emission from GRB 930131 and GRB 960924.

D.R. Willis¹,²,³ *, E.J. Barlow¹, A.J. Bird¹, D.J. Clark¹, A.J. Dean¹, M.L. McConnell⁴, L. Moran¹, S.E. Shaw¹,³, and V. Sguera¹

¹ School of Physics and Astronomy, University of Southampton, SO17 1BJ, UK
² Max-Planck-Institut fur extraterrestrische Physik, MPI, Garching, Munich, Germany
³ INTEGRAL Science Data Centre, CH-1290 Versoix, Switzerland
⁴ Space Science Center, University of New Hampshire, Durham, NH 03824, USA

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Abstract. The true nature of the progenitor to GRBs remains elusive; one characteristic that would constrain our understanding of the GRB mechanism considerably is gamma-ray polarimetry measurements of the initial burst flux. We present a method that interprets the prompt GRB flux as it Compton scatters off the Earth’s atmosphere, based on detailed modelling of both the Earth’s atmosphere and the orbiting detectors. The BATSE mission aboard the CGRO monitored the whole sky in the 20 keV – 1 MeV energy band continuously from April 1991 until June 2000. We present the BATSE Albedo Polarimetry System (BAPS), and show that GRB 930131 and GRB 960924 provide evidence of polarisation in their prompt flux that is consistent with degrees of polarisation of Π > 35% and Π > 50% respectively. While the evidence of polarisation is strong, the method is unable to strongly constrain the degree of polarisation beyond a systematics based estimation. Hence the implications on GRB theory are unclear, and further measurements essential.

Key words. gamma-rays: bursts – Techniques: polarimetric – Methods: data analysis – Polarization

1. Introduction

Whatever the mechanism, polarisation in the prompt gamma-ray flux from GRBs is evidence of strong magnetic fields within the burst. Theories on the GRB production mechanism can be constrained by different degrees of linear polarisation. For large degrees of polarisation, Π ≈ 80%, either shock accelerated synchrotron emission or a tuned Compton-drag model is the most likely [Lazzati et al. 2004; Coburn & Boggs 2003; Lyutikov et al. 2003]. For intermediate degrees of polarisation (20% < Π < 60%) two electromagnetic models have emerged that involve either synchrotron emission as the dominant source of radiation or as the result of viewing the burst from just outside the edge of the jet [Granot 2003; Ghisellini & Lazzati 1999]. Low degrees of polarisation can be a result of either a hydrodynamic model or from flux with a higher degree of polarisation experiencing partial depolarisation.

The high degree of linear polarisation initially reported in the prompt flux of GRB 021206 [degree of polarisation, Π = 80 ± 20%, Coburn & Boggs 2003] with the RHESSI experiment [McConnell et al. 2002] has stimulated much interest in the implications this has on GRB theory. Though GRB astronomy has the advantage of large fluxes, any polarimetric measurements are still dominated by systematic effects that can only be properly quantified by careful modelling. The importance of correctly evaluating the systematic effects is paramount in any measurement of GRB polarisation as was emphasised by Wigger et al. (2004) in the thorough re-analysis of the initial RHESSI result [Coburn & Boggs 2003]. Wigger et al. (2004) re-analysed the degree of polarisation to be Π = 41 ± 17%. This implies that many of the production mechanism theories are now similarly competitive.

Besides the RHESSI instrument [Lin et al. 2002], there are several other instruments that could be suitable gamma-ray polarimeters; IBIS and SPI aboard INTEGRAL [Winkler et al. 2003] have the as yet untested capability of detecting gamma-ray polarisation; COMPTEL [Schoenfelder et al. 1993] and BATSE, both aboard the Compton Gamma-Ray Observatory (CGRO) [Fishman et al. 1990] are also possible candidate instruments. For Compton polarimetry, the instrument is re-
quired to scatter the photons and detect the modulation induced by the polarisation-sensitive differential cross-section for Compton scattering (equation 1), see (Lei et al. 1994) for a review of Compton polarimetry in gamma-ray astronomy.

\[
\frac{d\sigma}{d\Omega} = \frac{r_0^2}{2} \left( \frac{E'}{E_0} \right)^2 \left( \frac{E'}{E_0} + \frac{E_0}{E'} - 2\sin^2\theta \cos^2\phi \right)
\]

(1)

where \(r_0^2\) is the classical radius of the electron, \(E_0\) the energy of the incident photon, \(E'\) the energy of the scattered photon, \(\theta\) the scatter angle and \(\phi\) the azimuthal angle relative to the polarisation direction. A problem in both COMPTEL and RHESSI is that each instrument has to both scatter and detect the linearly polarised photons. The scattering effective area, \(A_{\text{eff}}\), is small, resulting in a reduction in the number of double events required to detect the polarisation-induced azimuthal modulation. COMPTEL also had the problem that to act as a polarimeter the double-event rate had to be quite high (Lei et al. 1994), and at these high event rates vital kinematic information was lost in the telemetry. This is also expected to affect the polarimetry mode in IBIS to some extent.

To both scatter and detect the polarisation-induced azimuthal modulation, an instrument has to be either highly event selective and tracking sensitive or requires a large scattered flux, larger than the majority of GRBs can supply to RHESSI and INTEGRAL (Swinyard et al. 1994). The problem of tracking and selecting the events most efficiently is currently being addressed in missions such as the Tracking and Imaging Gamma-Ray Experiment (TIGRE) (O’Neill et al. 1996) and the Gamma-Ray Polarimeter Experiment (GRAPE) (McConnell et al. 2003). BATSE however is in the unique position to address the problem of increasing \(A_{\text{eff}}\) for scattering the incident gamma-ray photons. McConnell et al. (1994) proposed that polarimetry was possible by combining the large \(A_{\text{eff}}\) of the Earth’s atmosphere for Compton scattering with the 4\(\pi\) steradian field-of-view of the eight BATSE modules. Unfortunately the response of the CGRO to non-planar flux was difficult to handle analytically due to the uncollimated nature of BATSE and so could only be properly evaluated with modelling techniques. The modelling technique, called Mass Modelling, has been developed at Southampton and has been employed successfully on BATSE (Westmore 2002; Shaw et al. 2003, 2004 and references therein), INTEGRAL (Ferguson et al. 2003) and Swift (Willis 2002). The complex analytical situation and systematic effects resulting from inverting the CGRO in the GRB flux scattered off the Earth’s atmosphere can be modelled using this Mass Modelling approach, enabling polarimetry measurements with BATSE to take full advantage of the large \(A_{\text{eff}}\) that the Earth’s atmosphere provides.

Here we present the technique of albedo Compton polarimetry from the Earth’s Atmosphere (McConnell et al. 1994) combined with the BATSE Mass Model (BAMM) (Shaw et al. 2003), the BATSE Albedo Polarimetry System (BAPS).

2. The BATSE Albedo Polarimetry System (BAPS)

The BATSE mission aboard CGRO monitored the whole sky in the 20keV – 1MeV energy band continuously from April 1991 until June 2000. It was designed primarily to detect and locate GRBs (Fishman et al. 1995). BATSE consisted of eight uncollimated 2025cm\(^2\) NaI(Tl) Large Area Detectors (LADs), which had a combined field of view of 4\(\pi\) steradians. Also positioned below each LAD was a Spectral Detector (SD) made of 127 cm\(^2\) NaI(Tl). The SDs were optimized for energy coverage and resolution, and were operated over the range 10keV – 100MeV. An illustration of the CGRO and a BATSE module can be seen in Figure 1. BATSE was an extremely successful experiment and the current burst catalogue has over 8000 triggers (http://cossc.gsfc.nasa.gov/batse/BATSE_Ctlg/index.html).

The discovery that GRBs are located at cosmological distances is largely due to data collected with BATSE (Meegan et al. 1992).

BAPS uses both the BATSE Mass Model (BAMM) and a model of the Earth’s atmosphere. The system is highly CPU intensive and combining both models into one model is impractical due to the relative scales. The model of the Earth’s atmosphere was written in GEANT4 (Ravndal et al. 1993), in order to include the polarised Compton scattering process. It consists of a ball of radius 6503 km [Earth radius + 125 km] made of an integrated atmospheric composition of constant density \([\rho = 8.33 \times 10^{-5} \text{ g.cm}^{-3}, 38% \text{ O}_2 : 62% \text{ N}_2]\), the effects of this assumption are discussed in section 4.4.3. The BATSE volume was an imaginary sphere of radius 200 km at an altitude of 500 km. Figure 2 shows the scale of the Earth albedo model alongside a schematic of BAMM. BAMM (Shaw et al. 2003) was written in GEANT3 to dynamically flat-field the variable background for the purposes of occultation imaging (Harmon et al. 2002) and the subsequent production of an all-sky survey (Shaw et al. 2004). BAMM reproduces the instrument performance to a level of accuracy that is impossible through analytical techniques alone. For a full review of the mass modelling approach and its applications, see (Dean et al. 2003).

BAPS simulates the GRB flux as it scatters off the atmosphere; the distribution of this scattered flux is recorded as it passes through a volume equivalent to where the CGRO was at the time of the burst. Polarised flux preferentially scatters perpendicular to the direction of the polarisation vector, so any distribution produced as a result of polarised flux will appear as an anti-phased excess towards the limbs of the Earth, when viewed from the direction of the burst. The recorded distribution scattered from the atmosphere is then used as an input to the BAMM Monte Carlo and the count-rates simulated for each LAD. This process is carried out for an unpo-
polarised burst and for a fully polarised burst, for a range of polarisation angles. The simulated LAD count-rates are then compared to the burst count-rates in the BATSE data. The fitting procedure varies the percentage polarisation and the polarisation angle. The Z statistic, defined by

$$Z = \frac{R_1 - R_2}{\sqrt{R_1 t_1 + R_2 t_2}}$$

where $R$ and $t$ are the count-rates and observation times respectively, is used to compare the two count-rates: simulated and recorded.

BATSE uses two formats of data, CONTinuous (CONT) and DISCriminator Large Area (DISCLA). DISCLA data contains 4 broad energy channels with a binning width of 1.024 s. CONT data contains 16 energy channels with a binning width of 2.048 s. Due to BAPS only being sensitive to the Compton domain and below, CONT data is used to optimise this sensitivity. The sacrifice of better timing resolution is not a problem until time-evolving polarimetry is considered.

3. Burst Selection: GRB 930131 and GRB 960924

The selection criteria for a suitable burst are that the burst has to be short, strong, have a geocentre to burst angle, $\eta$, as close to 180° as possible (i.e. to have BATSE directly between the Earth and the burst) and for it to have no obvious evolving features. This selection reduces any background variations, maximise the polarisation-induced modulations and avoid the possibility of reprocessing at the source reducing the degree of polarisation, perhaps due to reverse shocks and scattering at the source. Suitable bursts were identified via a plot of burst strength against burst to geocentre angle, $\eta$ (Figure 3). Two of the strongest and most suitably placed BATSE triggers were identified as being suitable for this study: trigger 2151 (GRB 930131) and trigger 5614 (GRB 960924). The remaining triggers were either too long or had clear evolving features in the light curve, and so were discarded at this stage. The parameters of these two selected bursts are displayed in Table 1.

| GRB     | 930131 | 960924 |
|---------|--------|--------|
| Trigger | 2151   | 5614   |
| $\eta$  | 170.7  | 162.4  |
| CONT bins (2.048 seconds) | 2 | 4 |
| Band model $\alpha$ | -1.2 | -0.3 |
| Band model $\beta$ | -2.5 | -9.9 |
| Band model $E_{\text{break}}$ (keV) | 439 | 251 |

Fig. 3. All BATSE triggers plotted to identify the six most suitable candidates for polarimetry. Triggers 2151 and 5614 were initially selected with the remainder discarded due to their particularly long and structured lightcurves.
4. Analysis

4.1. Simulations

The first stage of BAPS is to produce the input distributions for the Mass Model. Fully polarised flux from the GRBs (with polarisation angle incremented every 15°), and an unpolarised burst flux were simulated incident on the atmosphere. 0.03% of the incident flux was seen to reflect and pass through the imaginary bubble containing the CGRO. The albedo angular distribution for polarised and unpolarised flux for both the GRBs can be seen in Figure 4.

BAMM was then used to simulate the effect these distributions of albedo flux have on the count-rates in the LADs during each burst. The direct GRB flux is also a simulated component. As seen in the albedo spectra (Figure 4) this technique is only sensitive in the Compton domain and so the energy range of the fitting is restricted to LAD CONT channels 3–6 (31–98 keV). The Compton domain extends up to higher energies but these restrictions equated to 67% and 62% of the albedo flux for GRB 930131 and GRB 960924 respectively.

4.2. LAD selection and fitting

The LADs selected are divided into two types, the first are used to normalise the model count-rates for the direct and albedo fluxes and the second are LADs pointed such that the albedo count-rate varies considerably depending on the degree and angle of polarisation. The remaining LADs are ignored if they have low simulated count-rates and low levels of polarisation-induced modulation, thus reducing the systematic and counting errors in the system and avoiding adding unconstrainable parameters into the fitting process. The component direct from the burst is normalised using the count-rate from the most burstward
The spectra of the albedo component of GRB 930131 and GRB 960924. The dotted lines represent the domain in which the data is compared: 31–98 keV (LAD CONT channels 3–6). These limits equate to 67% and 62% of the albedo flux respectively.

LAD. The simulated count-rates are fitted to the recorded count-rates with the normalisation of the direct flux fixed as free parameters. The ambient background is removed by taking a mean count-rate from before and after the burst. The fit statistic used is a simple Z-test for comparing the simulated and recorded count-rates.

The criteria for selecting the LADs is best demonstrated by looking at the simulated count-rate components for each LAD, see tables 2 and 3. The figure-of-merit consists of the percentage of simulated albedo flux, \( A \), in each LAD compared to the total flux, the maximum variation due to GRB polarisation angle and the level of albedo flux compared to the LAD with the maximum albedo flux. For each burst one LAD is selected to normalise the direct flux and two to measure polarisation-induced modulation.

### 4.2.1. GRB 930131

The LADs selected are LAD1, LAD2, LAD3, and LAD4. LAD1 and LAD2 have the largest albedo count-rate and are most sensitive to polarisation. LAD4 is used to set the normalisation for the direct flux and LAD3 is used to assist the normalisation of the albedo flux.

### 4.2.2. GRB 960924

The LADs selected are LAD0, LAD2, LAD4, and LAD6. LAD6 is used to set the normalisation for the direct flux. For this orientation of the CGRO, the polarisation-induced variation appears in all LADs. LAD0 is used to primarily assist the normalisation of the albedo flux whereas LAD2 and LAD4 show the largest albedo flux and have >30% variability due to polarisation angle and so are used to test any polarisation-induced modulation.

| LAD | \( \Delta A / A_{\text{Direct}} \) | \( \% \) | \( \Delta A_{\text{pol}} / \Delta A_{\text{unpol}} \) | \( \% \) | \( \Delta A / A_{\text{max}} \) | \( \% \) |
|-----|----------------------------------|------|---------------------------------|------|----------------------------|------|
| LAD0 | 1.0 | 45 | 38 |
| LAD1 | 0.49 | 38 | 59 |
| LAD2 | 0.70 | 17 | 60 |
| LAD3 | 0.73 | 7 | 100 |
| LAD4 | <1 | 58 | 1 |
| LAD5 | 0.16 | 25 | 2 |
| LAD6 | <1 | 38 | 9 |
| LAD7 | 0.09 | 31 | 17 |

| LAD | \( \Delta A / A_{\text{Direct}} \) | \( \% \) | \( \Delta A_{\text{pol}} / \Delta A_{\text{unpol}} \) | \( \% \) |
|-----|----------------------------------|------|---------------------------------|------|
| LAD0 | 31 | 11 | 100 |
| LAD1 | 51 | 10 | 29 |
| LAD2 | 8 | 30 | 57 |
| LAD3 | 7 | 35 | 14 |
| LAD4 | 3 | 36 | 62 |
| LAD5 | 4 | 44 | 13 |
| LAD6 | <1 | 21 | 8 |
| LAD7 | 4 | 41 | <1 |

### 4.3. Results

The selected LADs are simultaneously fitted to the CONT count-rates and the best fit for each incident polarisation angle is shown in Figures 6 and 8. The significance of each of these polarisation angle fits can be taken from Figures 7 and 9, the values for which were taken from a Z-test comparing the total count-rates in the LADs from the simulation and CONT data using the parameters supplied by the best fit for each polarisation angle. Also included in these plots of the Z statistic are the corresponding \( \Pi \) values.

GRB 930131 shows evidence of a high degree of polarisation. In this case, using simulated data with the polarisation angle at 90 degrees to the actual polarisation angle, the best fit will be to the unpolarised simulated flux, due to the phased nature of the excess at the Earth’s limb. Hence for this burst the degree of polarisation will only move towards a higher value when the fit is improved by the polarised GRB. Though this is seen in Figure 7 the sharp change in the degree of polarisation compared to the smoother change in the Z statistic suggests that this...
Fig. 6. A comparison of the fitted simulated count-rates with the LAD CONT data for GRB 930131. The CONT data is shown as a dashed line, the error on that count is shown as solid lines. The best fit for each polarisation angle is shown with the relevant counting error from the Monte Carlo. LAD3 and LAD4 simulations are expected to match the recorded data, independent of polarisation. LAD1 and LAD2 are sensitive to polarisation in the GRB's prompt flux and will show an improvement in fit for angles that corresponds to the burst polarisation angle.

Fig. 7. For GRB 930131, the Z statistic (solid line) for each of the best fits from Figure 6. The statistic becomes significant (95% confidence) between polarisation angles of 163° and 9° and this corresponds to a high (>90%) degree of polarisation (dashed line). The Earth’s North/South equivalent angle is at 5°, see section 4.4.1.

Fig. 9. For GRB 960924, the Z statistic (solid line) for each of the best fits from Figure 8. The statistic becomes significant (95% confidence) between polarisation angles of 58° and 64° and this corresponds to a high (100%) degree of polarisation (dashed line). The Earth’s North/South equivalent angle is at 47°, see section 4.4.1.

technique is not particularly sensitive to the degree of polarisation and can initially only be used as evidence of polarisation in the prompt GRB flux and that GRB 930131 is best fit by a fully polarised model (though see section 4.4.1 for an estimate of Π).
Fig. 8. A comparison of the fitted simulated count-rates with the LAD CONT data for GRB 960924. The CONT data is shown as a dashed line, the error on that count is shown as solid lines. The best fit for each polarisation angle is shown with the relevant counting error from the Monte Carlo. LAD6 simulations are expected to match the recorded data, independent of polarisation. LAD2 and LAD4 are sensitive to polarisation in the GRB's prompt flux and will show an improvement in fit for the specific angle that the burst is polarised to. LAD0 is largely insensitive to polarisation and so should be consistent with the CONT data regardless of polarisation.

GRB 960924 also shows evidence of a high degree of polarisation. The variation in the Z statistic in this case shows a clear dip between 35° and 120° and is flat for the remainder of angles. The sharp increase in the degree of polarisation compared to the dip in the Z statistic clearly shows the angles for which a polarised burst fits the CONT data better than an unpolarised burst. Again this is evidence of polarisation, that is best fit by a fully polarised model.

4.4. Systematic Effects

These results are expected to contain a series of systematic effects and a discussion and evaluation of these effects follows. The LAD selection by way of a figure-of-merit for each burst is expected to minimise these.

4.4.1. North-South Albedo Gamma-Ray Excess

The atmosphere itself is a strong source of gamma radiation. Cosmic ray interactions in the atmosphere radiate gamma-rays isotropically. This emissivity depends solely on the intensity of the incident cosmic rays and the atmospheric depth [Dean et al. 1989]. An angle-averaged spectrum of this albedo flux is also presented by Gehrels (1991). This albedo flux will be concentrated at the Earth’s limb as the largest portion of the atmosphere is being observed. Along the limb, the strongest source of albedo flux will be towards the polar regions, as the rigidity of the geomagnetic field deflecting the cosmic rays is at its weakest [Smart & Shea 1994]. This will result in an excess that will be reminiscent of the excess that BAPS will be looking for. During these bursts BATSE was at an inclination of -16° and -10° for GRB 930131 and GRB 960924 respectively. The emission from the south pole would have been visible to BATSE in both cases. The cosmic ray induced albedo flux would therefore falsely suggest a polarisation angle of 5° and 47° respectively, which corresponds to the inferred polarisation angle for GRB 930131 but not for GRB 960924. However, GRB 930131 should not be discounted as a possible source of prompt polarised flux as any excess produced from the South pole should be removed during the background subtraction within the fitting procedure.
4.4.2. East-West Cosmic Ray Anisotropy

Another effect of the geomagnetic field is the cosmic ray east-west anisotropy, from the direct detection of cosmic ray particles. The rigidity of the geomagnetic field indicates the ability of the field to fully deflect a charged particle. Faraday’s Law defines the direction that a cosmic ray is deflected in, and results in an excess in the westerly direction. This would correspond, in this case, to a falsely inferred polarisation angle of 95° and 137° for GRB 930131 and GRB 960924 respectively. These polarisation angles do not show an improvement in the fit and so this anisotropy is not responsible for a positive detection. Similar to the polar albedo excess discussed in section 4.4.1, any anisotropy should not be seen due to the background removal and it is not expected to change on the order of seconds. Moreover, any level of background due to cosmic ray interactions will be low as each BATSE LAD had a plastic scintillator vetoing cosmic ray events. Finally, it is estimated that both the North-South albedo excess and the East-West cosmic ray anisotropy are expected to be extended by 20–30°, leaving few of the potential polarisation angles free from these systematic effects. For weaker bursts, these North-South and East-West regions may give false detections.

4.4.3. Modelling Software

Combining two GEANT models will introduce a level of systematic error but due to the relative scale of the two models this is unavoidable. The GEANT 3 model of the CGRO is proven to work to an acceptable level of accuracy [Shaw et al. 2003]. However, the GEANT 4.6.1 model of the Earth’s atmosphere may introduce some problems. GEANT 4 is a new suite of software and is still under development, requiring verification at each stage. One area that requires modification was highlighted by Mizuno et al. (2004): the modulation produced from scattering a polarised beam. For their GEANT model the unmodified version of GEANT4 was seen to scatter a Crab (25–200keV) spectrum with around 30% variation with scattering angle, whereas the beam-test verified and modified version of GEANT was seen to give a modulation of around 64%. The models are entirely different but this suggests that the level of modulation in each LAD due to polarisation in BAPS could be up to half of what it should physically be. This is undoubtedly a systematic error which will degrade the sensitivity of BAPS but not one that would give a false indication of polarisation. Future runs will include the modifications suggested by Mizuno et al. (2004) and it is expected that BAPS will become far more sensitive as a result, and will give a better indication as to the degree of polarisation.

4.4.4. Constant Density Atmospheric Model

The model of the Earth’s Atmosphere consists of a ball of constant density (8.33 × 10^3 g.cm\(^{-3}\)) and composition (38% O\(_2\) : 62% N\(_2\)). This is an approximation that will introduce a systematic error but aid in the simplicity and speed of BAPS. The constant density was achieved by averaging over the first 125 km of the atmosphere. To evaluate the effectiveness of this approximation, a single column element of this atmospheric model was compared to a graded (in density and composition) 12 layered model of the atmosphere. The rate of back-scattered events (scatter angle >90°) and resultant direction of the photon emerging from the atmosphere were recorded for both models, as seen in Figure 10. For the layered model, the majority of the scattering occurs in the 20–50 km region, and the 0–125 km integration reproduces 99.7% of these back-scattered events. The back-scattered zenith angle distribution for an on-axis beam of photons and one at 40° off-axis was also compared for the two models. There is a slight excess around the scatter angle of 160° and below for the approximated model. However, it is clear that this discrepancy between the two models will not effect the BAPS fitting when considering the uncollimated nature of the BATSE LADs. Furthermore, the 0.3% change in the albedo count-rate will be accounted for in the fitting routine. The diametric excess that polarisation will induce will not be effected.

![Fig. 10. A comparison between using a 12 layered graded atmospheric model and a constant density and composition integrated model. The incident photon energy is 200keV. The upper plot shows the back-scatter rate with altitude and the lower two histograms show the albedo zenith distribution.](image-url)
This depth-dependent scattering rate also discards the need for including the Earth itself into the model. A photon interacting such that it scatters in the atmosphere towards the direction of the CGRO but must travel through the Earth to reach the CGRO volume will certainly be absorbed by the atmosphere in the model, and so discarded.

4.5. Estimating the Degree of Polarisation

The physical implications of polarised prompt flux from GRBs hinges on Π, the percentage of the flux that is polarised. As BAPS normalises the levels of albedo flux and involves comparisons of count-rates, the degree to which these GRBs are polarised cannot be measured directly. However, an estimate of Π can be achieved through restricting Π to a maximum value and inspecting the change in the Z fitting statistic. The data is refitted with restrictions applied to Π, as shown in Figures 11 and 12.

![Fig. 11. The CONT data for GRB 930131 is compared to the simulations with Π restricted to a range of maximum values between 0% and 100%. The confidence levels in the Z statistic are shown.](image)

![Fig. 12. The CONT data for GRB 960924 is compared to the simulations with Π restricted to a range of maximum values between 0% and 100%. The confidence levels in the Z statistic are shown.](image)

5. Conclusions

BAPS has been presented here in the context of GRB 930131 and GRB 960924, both of which show a high degree of polarisation. The caveats to this result have been discussed and no systematic explanation can be found for these positive detections of polarisation. The GRB selection procedure is initially highly restrictive, however there are several bursts that are also suitable to be analysed by BAPS in future studies. The initial stage of the BAPS analysis (modelling of the albedo flux) was run for a variety of angles taken from figure 3 and it was seen that the difference between the polarised and unpolarised cases was statistically significant up to \( \eta < 120^\circ \), suggesting that BAPS is primarily restricted in its sensitivity by the number of LADs that are illuminated by the GRB albedo flux and not the off-axis angle of the burst. This increases the number of GRBs that can be analysed with BAPS, however since the sensitivity is governed by the number of illuminated LADs, it removes the possibility of analysing the systematics of a suitably off-axis GRB that is unable to exhibit any polarisation effects.

Though the significance of the fit for GRB 930131 decreases as the maximum allowed value for Π decreases, the fit is still significant down to a Π values of around 70%–80%. However, the fit for GRB 960924 suffers more severely with any change to Π and the significance of the result is lost at a level of 90% polarisation. This value for Π is unphysical and so the systematics discussed in section 4.4.3 must be considered. As it is quite possible that the polarisation-induced modulation is underestimated by up to a factor of 2, the value of Π may well be overestimated by a similar figure. Therefore the degree of polarisation in GRB 930131 is consistent with Π ranging between 35% and 100%; and GRB 960624 is consistent with Π ranging between 50% and 100%.

There are areas in BAPS that would benefit from further augmentation: modification to the current GEANT4 release, an analytical model of the scattering off the Earth’s atmosphere and a full analysis of the systematic errors to enable a quantitative estimate to the degree of polarisation and enable the inclusion of more LADs into the fitting procedure.

Though the degree of polarisation is an estimate from these results, there is distinct evidence that polarisation exists in prompt GRB flux and to a degree that is large enough to be detected by viewing the Earth’s atmosphere with the BATSE detectors. The implications of this on the GRB mechanism are that a large scale magnetic field is present in the emitting region, conclusions beyond this will remain uncertain until future measurements can be made.
Viewing the prompt flux is crucial to future polarimetry measurements, something that is only currently possible through a large scale survey of the BATSE catalogue with BAPS or through an on-axis INTEGRAL burst such as the recent strong burst GRB 041219.

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