Semilinear wave equation on compact Lie groups

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Abstract
In this note, we study the semilinear wave equation with power nonlinearity $|u|^p$ on compact Lie groups. First, we prove a local in time existence result in the energy space via Fourier analysis on compact Lie groups. Then, we prove a blow-up result for the semilinear Cauchy problem for any $p > 1$, under suitable sign assumptions for the initial data. Furthermore, sharp lifespan estimates for local (in time) solutions are derived.

Keywords Blow-up · Wave equation · Lifespan estimates · Compact Lie group · Local existence

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1 Introduction

Let $\mathbb{G}$ be a compact Lie group and let $\mathcal{L}$ be the Laplace–Beltrami operator on $\mathbb{G}$ (which coincides with the Casimir element of the enveloping algebra). In the present work, we prove a blow-up result for the Cauchy problem for the semilinear wave equation with power nonlinearity, namely,

$$\begin{cases}
\partial_t^2 u - \mathcal{L} u = |u|^p, & x \in \mathbb{G}, \ t > 0, \\
u(0, x) = \varepsilon u_0(x), & x \in \mathbb{G}, \\
\partial_t u(0, x) = \varepsilon u_1(x), & x \in \mathbb{G},
\end{cases}$$

where $p > 1$ and $\varepsilon$ is a positive constant describing the smallness of Cauchy data.

Throughout the paper $L^q(\mathbb{G})$ denotes the space of $q$-summable functions on $\mathbb{G}$ with respect to the normalized Haar measure for $1 \leq q < \infty$ (respectively, essentially
bounded for $q = \infty$) and for $s > 0$ and $q \in (1, \infty)$ the Sobolev space $H^{s,q}_L(G)$ is defined as the space

$$H^{s,q}_L(G) = \left\{ f \in L^q(G) : (-L)^{s/2} f \in L^q(G) \right\}$$

dowed with the norm $\|f\|_{H^{s,q}_L(G)} = \|f\|_{L^q(G)} + \|(-L)^{s/2} f\|_{L^q(G)}$. As customary, the Hilbert space $H^{s,2}_L(G)$ will be simply denoted by $H^s_L(G)$.

For the classical semilinear wave equation in $\mathbb{R}^n$ it has been proved that the critical exponent is the positive root $p_{\text{Str}}(n)$ of the quadratic equation

$$(n-1)p^2 - (n+1)p - 2 = 0,$$

which is the so-called Strauss exponent, named after the author of [16]. For a detailed overview and a complete list of references on the proof of the Strauss’ conjecture we refer to the introduction of [18].

In the not-flat case, the semilinear wave equation has been studied in many Lorentzian metrics such as Schwarzschild spacetime [1,8], de Sitter spacetime [3,5,20–22] and Einstein–de Sitter spacetime [6,11,12,19]. In the present work, we will consider the case of semilinear wave equations in compact Lie groups, which seems to be an almost unexplored topic in the literature, to the knowledge of the author.

By employing Fourier analysis for compact Lie groups, the local well-posedness of the semilinear Cauchy problem (1) in the energy evolution space, that is $C([0, T], H^1_L(G)) \cap C^1([0, T], L^2(G))$, will be proved. In particular, a Gagliardo–Nirenberg type inequality recently proved in [15] will be used in order to estimate the power nonlinearity in $L^2(G)$. As byproduct of this local existence result we will obtain lower bound estimates for the lifespan of a local in time solution (that is, the maximal existence time of a solution).

Finally, we will prove the nonexistence of globally in time defined solutions to (1) for any $p > 1$ and regardless of the size of initial data, provided that the data fulfill suitable sign assumptions.

1.1 Main results

Let us get started with the statement of the local existence result for the semilinear Cauchy problem (1).

**Theorem 1** Let $G$ be a compact, connected Lie group. Let us assume that the topological dimension $n$ of $G$ satisfies $n \geq 3$. Let $(u_0, u_1) \in H^1_L(G) \times L^2(G)$ and let $p > 1$ such that $p \leq \frac{n}{n-2}$. Then, there exists $T = T(\varepsilon) > 0$ such that the Cauchy problem (1) admits a uniquely determined mild solution

$$u \in C \left([0, T], H^1_L(G)\right) \cap C^1 \left([0, T], L^2(G)\right).$$
Furthermore, the lifespan $T$ satisfies the following lower bound estimates

$$T(\varepsilon) \geq \begin{cases} \varepsilon^{1-rac{n}{p+1}} & \text{if } u_1 \neq 0, \\ \varepsilon^{1-rac{n-1}{2}} & \text{if } u_1 = 0, \end{cases}$$

where the positive constant $c$ is independent of $\varepsilon$.

**Remark 1** The upper bound assumption for the exponent $p$ in Theorem 1 is made in order to apply a Gagliardo–Nirenberg type inequality proved in [15] [Remark 1.7]. The restriction $n \geq 3$ is made to fulfill the assumptions for the employment of such inequality as well and it can be avoided by looking for solutions in weaker spaces than the one in the statement of Theorem 1, i.e., in $C \left([0, T], H_s^1(G) \cap C^1([0, T], L^2(G))\right)$ for some $s \in (0, 1)$.

We introduce now a suitable notion of energy solutions for the semilinear Cauchy problem (1).

**Definition 1** Let $(u_0, u_1) \in H_s^1(G) \times L^2(G)$. We say that

$$u \in C \left([0, T], H_s^1(G)\right) \cap C^1 \left([0, T], L^2(G)\right) \cap L_p^p([0, T] \times G)$$

is an energy solution on $[0, T)$ to (1) if $u$ fulfills the integral relation

$$\begin{align*}
\int_G \partial_t u(t, x) \psi(t, x) \, dx - \int_G u(t, x) \psi_x(t, x) \, dx &= -\varepsilon \int_G u_1(x) \psi(0, x) \, dx \\
+ \varepsilon \int_G u_0(x) \psi_x(0, x) \, dx + \int_0^t \int_G u(s, x) \left(\psi_{ss}(s, x) - L\psi(s, x)\right) \, dx \, ds \\
&= \int_0^t \int_G |u(s, x)|^p \psi(s, x) \, dx \, ds \tag{3}
\end{align*}$$

for any $\psi \in C_0^\infty([0, T] \times G)$ and any $t \in (0, T)$.

**Theorem 2** Let $G$ be a compact Lie group. Let $(u_0, u_1) \in H_s^1(G) \times L^2(G)$ be non-negative and not both trivial functions. Let $p > 1$ and let

$$u \in C \left([0, T], H_s^1(G)\right) \cap C^1 \left([0, T], L^2(G)\right) \cap L_p^p([0, T] \times G)$$

be an energy solution to (1) according to Definition 1 with lifespan $T = T(\varepsilon)$. Then, there exists a positive constant $\varepsilon_0 = \varepsilon_0(u_0, u_1, p) > 0$ such that for any $\varepsilon \in (0, \varepsilon_0]$ the energy solution $u$ blows up in finite time. Moreover, the upper bound estimates for the lifespan

$$T(\varepsilon) \leq \begin{cases} C\varepsilon^{1-rac{n}{p+1}} & \text{if } u_1 \neq 0, \\ C\varepsilon^{1-rac{n-1}{2}} & \text{if } u_1 = 0, \end{cases}$$

hold, where the constant $C > 0$ is independent of $\varepsilon$. 

Remark 2 Combining (2) and (4), we obtain the sharp lifespan estimates

\[
\begin{align*}
&c \varepsilon^{-\frac{p-1}{p+1}} \leq T(\varepsilon) \leq C \varepsilon^{-\frac{p-1}{p+1}} \text{ if } u_1 \neq 0, \\
&c \varepsilon^{-\frac{p-1}{2}} \leq T(\varepsilon) \leq C \varepsilon^{-\frac{p-1}{2}} \text{ if } u_1 = 0,
\end{align*}
\]

for local in time solutions to (1). Therefore, the nontriviality of \( u_1 \) plays a crucial role on the lifespan estimates, analogously to what happens for the Euclidean case (cf. [17]).

Notations

Hereafter we use the following notations: \( \mathcal{L} \) denotes the Laplace–Beltrami operator on \( G \); \( \text{Tr}(A) = \sum_{j=1}^{d} a_{jj} \) and \( A^* = (a_{ji})_{1 \leq i, j \leq d} \) denote the trace and the adjoint matrix of \( A = (a_{ij})_{1 \leq i, j \leq d} \in \mathbb{C}^{d \times d} \), respectively; \( I_d \in \mathbb{C}^{d \times d} \) denotes the identity matrix; \( dx \) stands for the normalized Haar measure on the compact group \( G \); finally, when there exists a positive constant \( C \) such that \( f \leq C g \) we write \( f \lesssim g \).

2 Local existence in energy space

In this section, we will prove Theorem 1. First, we recall the notion of mild solutions to (1). By Duhamel’s principle, the solution to the linear problem

\[
\begin{align*}
\partial_t^2 u - \mathcal{L} u &= F(t, x), \quad x \in G, \quad t > 0, \\
u(0, x) &= u_0(x), \quad x \in G, \\
\partial_t u(0, x) &= u_1(x), \quad x \in G
\end{align*}
\]

(5)
can be represented as

\[
u(t, x) = u_0(x) *_{(x)} E_0(t, x) + u_1(x) *_{(x)} E_1(t, x)
+ \int_0^t F(s, x) *_{(x)} E_1(t - s, x) \, ds,
\]

where \( E_0(t, x) \) and \( E_1(t, x) \) denote, respectively, the fundamental solutions to (5) in the homogeneous case \( F = 0 \) with initial data \( (u_0, u_1) = (\delta_0, 0) \) and \( (u_0, u_1) = (0, \delta_0) \). Let us point out that, in order to get the previous representation formula, we applied the invariance by time translations for the wave operator \( \partial_t^2 - \mathcal{L} \) and the property \( L(v *_{(x)} E_1(t, \cdot)) = v *_{(x)} L(E_1(t, \cdot)) \) for any left-invariant differential operator \( L \) on \( G \).
Also, \( u \) is said a **mild solution** to (1) on \([0, T]\) if \( u \) is a fixed point for the nonlinear integral operator \( N \) defined by

\[
N : u \in X(T) \rightarrow Nu(t, x) = \varepsilon u_0(x) *_{(x)} E_0(t, x) + \varepsilon u_1(x) *_{(x)} E_1(t, x) + \int_0^t |u(s, x)|^p *_{(x)} E_1(t - s, x) \, ds
\]

on the evolution space \( X(T) \doteq C \left([0, T], H^1_L(G) \right) \cap C^1 \left([0, T], L^2(G) \right), \)

endowed with the norm

\[
\|u\|_{X(T)} = \sup_{t \in [0, T]} \left( \alpha(t)^{-1}\|u(t, \cdot)\|_{L^2(G)} + \|(-\mathcal{L})^{1/2}u(t, \cdot)\|_{L^2(G)} + \|\partial_t u(t, \cdot)\|_{L^2(G)} \right),
\]

where

\[
\alpha(t) = \begin{cases} 
(1 + t) & \text{if } u_1 \neq 0; \\
1 & \text{if } u_1 = 0.
\end{cases}
\]

As we will see in Proposition 1, the function \( \alpha \) represents the long time behavior of the \( L^2(G) \) norm of the solution to the corresponding linear homogeneous Cauchy problem. The introduction of this time-dependent weight in the norm on \( X(T) \) will allow us to determine sharp lower bound estimates for the lifespan of local in time solutions to (1).

By employing Banach’s fixed point theorem, we will show that \( N \) admits a uniquely determined fixed point and we will provide a lower bound estimate for the lifespan \( T = T(\varepsilon) > 0 \). However, before considering the semilinear Cauchy problem, we determine \( L^2(G) - L^2(G) \) estimates for the solution to the corresponding linear homogeneous problem via the group Fourier transform with respect to the spatial variable. After that these estimates will have been obtained, we could show the local well-posedness for (1) by applying a Gagliardo–Nirenberg type inequality derived recently in [15] (cf. Lemma 1). The remaining part of this section is organized in the following way: in Sect. 2.1 we recall the main tools from Fourier Analysis on compact Lie groups which are necessary for our approach; hence, in Sect. 2.2 \( L^2(G) - L^2(G) \) estimates for the solution of the corresponding homogeneous linear problem and its first order derivatives are derived; finally, in Sect. 2.3 it will be shown that the operator \( N \) is a contraction on \( X(T) \).

### 2.1 Group Fourier transform

Let us recall some results on Fourier Analysis on compact Lie groups as in [9][Section 2.1]. For further details on this topic we refer to the monographs [4, 13].

A **continuous unitary representation** \( \xi : G \rightarrow \mathbb{C}^{d_{\xi} \times d_{\xi}} \) of dimension \( d_{\xi} \) is a continuous group homomorphism from \( G \) to the group of unitary matrix \( U(d_{\xi}, \mathbb{C}) \), that is, \( \xi(xy) = \xi(x)\xi(y) \) and \( \xi(x)^* = \xi(x)^{-1} \) for all \( x, y \in G \) and the elements \( \xi_{ij} : G \rightarrow \mathbb{C} \) of the matrix representation \( \xi \) are continuous functions for all \( i, j \in \{1, \ldots, d_{\xi}\} \). Two
representations $\xi, \eta$ of $G$ are said equivalent if there exists an invertible intertwining operator $A$ such that $A \xi(x) = \eta(x) A$ for any $x \in G$. A subspace $W \subset \mathbb{C}^{d_c}$ is $\xi$—invariant if $\xi(x) \cdot W \subset W$ for any $x \in G$. A representation $\xi$ is irreducible if the only $\xi$—invariant subspaces are the trivial ones $\{0\}, \mathbb{C}^{d_c}$.

The unitary dual $\hat{G}$ of the compact Lie group $G$ consists of the equivalence class $[\xi]$ of continuous irreducible unitary representation $\xi : G \to \mathbb{C}^{d_c} \times \mathbb{C}^{d_c}$.

Given $f \in L^1(G)$, its Fourier coefficients at $[\xi] \in \hat{G}$ is defined by

$$\hat{f}(\xi) = \int_G f(x) \xi(x)^* \, dx \in \mathbb{C}^{d_c} \times \mathbb{C}^{d_c},$$

where the integral is taken with respect to the Haar measure on $G$.

If $f \in L^2(G)$, then, the Fourier series representation for $f$ is given by

$$f(x) = \sum_{[\xi] \in \hat{G}} d_c \, \text{Tr} \left( \xi(x) \hat{f} \right),$$

where hereafter just one irreducible unitary matrix representation is picked in the sum for each equivalence class $[\xi]$ in $\hat{G}$. Furthermore, for $f \in L^2(G)$ Plancherel formula takes the following form

$$\|f\|_{L^2(G)}^2 = \sum_{[\xi] \in \hat{G}} d_c \|\hat{f}(\xi)\|_{\text{HS}}^2,$$

(6)

where the Hilbert–Schmidt norm of the matrix $\hat{f}(\xi)$ is defined as follows:

$$\|\hat{f}(\xi)\|_{\text{HS}}^2 = \text{Tr} \left( \hat{f}(\xi) \hat{f}(\xi)^* \right) = \sum_{i,j=1}^{d_c} \left| \hat{f}(\xi)_{ij} \right|^2.$$

For our analysis it is very important to understand the behavior of the group Fourier transform with respect to the Laplace–Beltrami operator $L$. Given $[\xi] \in \hat{G}$, then, all $\xi_{ij}$ are eigenfunctions for $L$ with the same not positive eigenvalue $-\lambda^2_\xi$, namely,

$$-L \xi_{ij}(x) = \lambda^2_\xi \xi_{ij}(x) \quad \text{for any } x \in G \text{ and for all } i, j \in \{1, \ldots, d_c\}.$$

This means that the symbol of $L$ is

$$\sigma_L(\xi) = -\lambda^2_\xi I_{d_c},$$

(7)

that is, $\hat{L} \hat{f}(\xi) = \sigma_L(\xi) \hat{f}(\xi) = -\lambda^2_\xi \hat{f}(\xi)$ for any $[\xi] \in \hat{G}$.

Finally, through Plancherel formula for $s > 0$ we have

$$\|f\|_{H^s_L(G)}^2 = \|(-L)^{s/2} f\|_{L^2(G)}^2 = \sum_{[\xi] \in \hat{G}} d_c \lambda^2_\xi \|\hat{f}(\xi)\|_{\text{HS}}^2.$$
2.2 $L^2(G)$–$L^2(G)$ estimates for the solution to the linear homogeneous problem

In this section, we derive $L^2(G)$–$L^2(G)$ estimates for the solution to the homogeneous problem associated with (1), that is,

$$
\begin{align*}
\partial_t^2 u - Lu &= 0, & x \in G, & t > 0, \\
u(0, x) &= u_0(x), & x \in G, \\
\partial_t u(0, x) &= u_1(x), & x \in G.
\end{align*}
$$

We follow the main ideas from [7,9], so, the group Fourier transform with respect to the spatial variable $x$ is applied together with Plancherel identity to determine an explicit expression for the $L^2(G)$ norms of $u(t, \cdot)$, $(\partial_t - L)^{1/2}u(t, \cdot)$ and $\partial_t u(t, \cdot)$, respectively.

Let $u$ be a solution to (8). Let $\hat{u}(t, \xi) = (\hat{u}(t, \xi)_{k\ell})_{1 \leq k, \ell \leq d}$, $[\xi] \in \hat{G}$ denote the group Fourier transform of $u$ with respect to the $x$—variable. Also, $\hat{u}(t, \xi)$ is a solution of the Cauchy problem for the system of ODEs (with size depending on the representation $[\xi]$)

$$
\begin{align*}
\partial_t^2 \hat{u}(t, \xi) - \sigma_{[\xi]}(\xi)\hat{u}(t, \xi) &= 0, & t > 0, \\
\hat{u}(0, \xi) &= \hat{u}_0(\xi), \\
\partial_t \hat{u}(0, \xi) &= \hat{u}_1(\xi).
\end{align*}
$$

By (7), it follows that the previous system is decoupled in $d^2$ independent scalar ODEs, namely,

$$
\begin{align*}
\partial_t^2 \hat{u}(t, \xi)_{k\ell} + \lambda_{[\xi]}^2 \hat{u}(t, \xi)_{k\ell} &= 0, & t > 0, \\
\hat{u}(0, \xi)_{k\ell} &= \hat{u}_0(\xi)_{k\ell}, \\
\partial_t \hat{u}(0, \xi)_{k\ell} &= \hat{u}_1(\xi)_{k\ell},
\end{align*}
$$

for any $k, \ell \in \{1, \ldots, d\}$. Straightforward computations provide the following representation formula for the solution to the parameter dependent ordinary linear homogeneous Cauchy problem (9)

$$
\hat{u}(t, \xi)_{k\ell} = G_0(t, \xi)\hat{u}_0(\xi)_{k\ell} + G_1(t, \xi)\hat{u}_1(\xi)_{k\ell},
$$

where

$$
G_0(t, \xi) = \begin{cases} 
\cos(\lambda_{[\xi]}t) & \text{if } \lambda_{[\xi]}^2 > 0, \\
1 & \text{if } \lambda_{[\xi]}^2 = 0,
\end{cases} \quad G_1(t, \xi) = \begin{cases} 
\sin(\lambda_{[\xi]}t) \frac{1}{\lambda_{[\xi]}} & \text{if } \lambda_{[\xi]}^2 > 0, \\
t & \text{if } \lambda_{[\xi]}^2 = 0.
\end{cases}
$$

Notice that $G_0(t, \xi) = \partial_t G_1(t, \xi)$ for any $[\xi] \in \hat{G}$. We underline explicitly that $0$ is an eigenvalue for the continuous irreducible unitary representation $1 : x \in G \to 1 \in \mathbb{C}$.  

Estimate for $\|u(t)\|_{L^2(G)}$

By (10) it follows

$$|\hat{u}(t, \xi)_{k\ell}| \leq |\hat{u}_0(\xi)_{k\ell}| + t |\hat{u}_1(\xi)_{k\ell}| \quad \text{for any } t \geq 0.$$  \hfill (12)

Consequently, by using Plancherel formula twice, we get

$$\|u(t, \cdot)\|^2_{L^2(G)} = \sum_{[\xi] \in \hat{G}} d_{\xi} \sum_{k, \ell=1}^{d_{\xi}} |\hat{u}(t, \xi)_{k\ell}|^2 \lesssim \sum_{[\xi] \in \hat{G}} d_{\xi} \sum_{k, \ell=1}^{d_{\xi}} \left( |\hat{u}_0(\xi)_{k\ell}|^2 + t^2 |\hat{u}_1(\xi)_{k\ell}|^2 \right)$$

$$= \|u_0\|^2_{L^2(G)} + t^2 \|u_1\|^2_{L^2(G)}. \hfill (13)$$

Estimate for $\|(-\mathcal{L})^{1/2} u(t)\|_{L^2(G)}$

By Plancherel formula, we find

$$\|(-\mathcal{L})^{1/2} u(t, \cdot)\|^2_{L^2(G)} = \sum_{[\xi] \in \hat{G}} d_{\xi} \|\sigma(-\mathcal{L})^{1/2}(\xi) \hat{u}(t, \xi)\|^2_{HS} = \sum_{[\xi] \in \hat{G}} d_{\xi} \sum_{k, \ell=1}^{d_{\xi}} \lambda_{\xi}^2 |\hat{u}(t, \xi)_{k\ell}|^2.$$

Since

$$\lambda_{\xi} |\hat{u}(t, \xi)_{k\ell}| \leq \lambda_{\xi} |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}|, \hfill (14)$$

then, from the previous identity we have

$$\|(-\mathcal{L})^{1/2} u(t, \cdot)\|^2_{L^2(G)} \lesssim \sum_{[\xi] \in \hat{G}} d_{\xi} \sum_{k, \ell=1}^{d_{\xi}} \left( \lambda_{\xi}^2 |\hat{u}_0(\xi)_{k\ell}|^2 + |\hat{u}_1(\xi)_{k\ell}|^2 \right)$$

$$= \|u_0\|^2_{H^1_G} + \|u_1\|^2_{L^2(G)},$$

where in the last step we applied again Plancherel formula.

Estimate for $\|\partial_t u(t)\|_{L^2(G)}$

Elementary computations show that for any $[\xi] \in \hat{G}$ and any $k, \ell \in \{1, \ldots, d_{\xi}\}$ it holds

$$\partial_t \hat{u}(t, \xi)_{k\ell} = -\lambda_{\xi}^2 G_1(t, \xi) \hat{u}_0(\xi)_{k\ell} + G_0(t, \xi) \hat{u}_1(\xi)_{k\ell},$$

where $G_0(t, \xi), G_1(t, \xi)$ are defined in (11). Therefore,

$$|\partial_t \hat{u}(t, \xi)_{k\ell}| \leq \lambda_{\xi} |\hat{u}_0(\xi)_{k\ell}| + |\hat{u}_1(\xi)_{k\ell}|.$$
Since the right-hand side of the previous inequality is the same one as in (14), we get immediately
\[ \| \partial_t u(t, \cdot) \|_{L^2(G)}^2 \lesssim \| u_0 \|_{H^1_L(G)}^2 + \| u_1 \|_{L^2(G)}^2. \]

Summarizing, in this section we proved the following proposition.

**Proposition 1** Let us assume \((u_0, u_1) \in H^1_L(G) \times L^2(G)\) and let
\[ u \in C([0, \infty), H^1_L(G)) \cap C^1([0, \infty), L^2(G)) \]
be the solution to the homogeneous Cauchy problem (8). Then, the following \(L^2(G)\)–\(L^2(G)\) estimates are satisfied
\[ \| u(t, \cdot) \|_{L^2(G)} \leq C \left( \| u_0 \|_{L^2(G)} + t \| u_1 \|_{L^2(G)} \right), \]
\[ \| (-L)^{1/2} u(t, \cdot) \|_{L^2(G)} \leq C \left( \| u_0 \|_{H^1_L(G)} + \| u_1 \|_{L^2(G)} \right), \]
\[ \| \partial_t u(t, \cdot) \|_{L^2(G)} \leq C \left( \| u_0 \|_{H^1_L(G)} + \| u_1 \|_{L^2(G)} \right), \]
for any \( t \geq 0 \), where \( C \) is a positive multiplicative constant.

### 2.3 Proof of Theorem 1

A fundamental tool to prove the local existence result is the following Gagliardo–Nirenberg type inequality, whose proof can be found in [15] (see also [9, Corollary 2.3]).

**Lemma 1** Let \( G \) be a connected unimodular Lie group with topological dimension \( n \geq 3 \). For any \( q \geq 2 \) such that \( q \leq \frac{2n}{n-2} \) the following Gagliardo–Nirenberg type inequality holds
\[ \| f \|_{L^q(G)} \lesssim \| f \|_{H^1_L(G)}^{\theta(n,q)} \| f \|_{L^2(G)}^{1-\theta(n,q)} \]
for any \( f \in H^1_L(G) \), where \( \theta(n, q) = n \left( \frac{1}{2} - \frac{1}{q} \right) \).

We can now proceed with the proof of Theorem 1. Let us estimate \( \| Nu \|_{X(T)} \) for \( u \in X(T) \). We rewrite \( Nu = u^{ln} + Ju \), where
\[ u^{ln}(t, x) \doteq \varepsilon u_0(x) *_{(x)} E_0(t, x) + \varepsilon u_1(x) *_{(x)} E_1(t, x), \]
\[ Ju(t, x) \doteq \int_0^t |u(s, x)|^p *_{(x)} E_1(t-s, x) \, ds. \]
By Proposition 1 we get immediately \( \|u^n\|_{X(T)} \lesssim \varepsilon \|(u_0, u_1)\|_{H^1_\omega(G) \times L^2(G)} \). On the other hand, due to the invariance by time translations of (8), it results

\[
\|a^j_t (-L)^{j/2} Ju(t, \cdot)\|_{L^2(G)} \lesssim \int_0^t (t-s)^{j-1+i} \|u(s, \cdot)\|^p_{L^2_p(G)} \, ds \\
\lesssim \int_0^t (t-s)^{j-1+i} \|u(s, \cdot)\|^p_{L^2_p(G)} \|u(s, \cdot)\|_{L^2(G)}^{p\beta(n,2,p)} \, ds \\
\lesssim \int_0^t (t-s)^{j-1+i} a(s)^p \|u\|_{X(t)}^p \, ds \\
\lesssim t^{2-(j+i)} a(t)^p \|u\|_{X(t)}^p
\]

(19)

for \( i, j \in \{0, 1\} \) such that \( 0 \leq i + j \leq 1 \). We stress that employing (18) in the above estimate, we have to require the condition \( p \leq \frac{n}{n-2} \) on \( p \) in Theorem 1. Similarly, combining Hölder’s inequality and (18), for \( i, j \in \{0, 1\} \) such that \( 0 \leq i + j \leq 1 \) we find

\[
\|a^j_t (-L)^{j/2} (Ju(t, \cdot) - Jv(t, \cdot))\|_{L^2(G)} \\
\lesssim \int_0^t (t-s)^{j-1+i} \|u(s, \cdot)\|^p - \|v(s, \cdot)\|^p_{L^2(G)} \, ds \\
\lesssim \int_0^t (t-s)^{j-1+i} \|u(s, \cdot)\|_{L^2_p(G)} \|u(s, \cdot)\|_{L^2(G)}^{p-1} \|v(s, \cdot)\|_{L^2_p(G)}^{p-1} \, ds \\
\lesssim t^{2-(j+i)} a(t)^p \|u - v\|_{X(t)} \left( \|u\|_{X(t)}^{p-1} + \|v\|_{X(t)}^{p-1} \right).
\]

(20)

\[
\|Nu\|_{X(T)} \leq \tilde{C} \varepsilon \|(u_0, u_1)\|_{H^1_\omega(G) \times L^2(G)} + \tilde{C} (1 + t)^{\beta(p)} \|u\|_{X(t)}^p, \\
\|Nu - Nv\|_{X(T)} \leq \tilde{C} (1 + t)^{\beta(p)} \|u - v\|_{X(t)} \left( \|u\|_{X(T)}^{p-1} + \|v\|_{X(T)}^{p-1} \right),
\]

where

\[
\beta(p) = \begin{cases} 
p + 1 & \text{if } u_1 \neq 0, \\
2 & \text{if } u_1 = 0.
\end{cases}
\]

Consequently, denoting \( R_0 = \|(u_0, u_1)\|_{H^1_\omega(G) \times L^2(G)} \), we find that for any \( R \geq 2\tilde{C} R_0 \) and any \( t \leq C \varepsilon^{-\frac{p-1}{\beta(p)}} \), where \( C \doteq (4\tilde{C} R)^{-\frac{1}{\beta(p)}} \), the mapping \( N \) satisfies

\[
\|Nu\|_{X(t)} \leq R \varepsilon, \quad \|Nu - Nv\|_{X(t)} \leq \frac{1}{2} \|u - v\|_{X(t)},
\]

for all \( u, v \in \mathcal{B}(R \varepsilon) \doteq \{ u \in X(t) : \|u\|_{X(t)} \leq R \varepsilon \} \), that is, \( N \) is a contraction mapping on the ball \( \mathcal{B}(R \varepsilon) \) in the Banach space \( X(t) \). Thus, Banach’s fixed point implies the existence of a uniquely determined fixed point \( u \) for \( N \), which is the mild solution to (1) on \([0, t] \subset [0, T(\varepsilon)]\) we were looking for. Besides, we got the lower bound estimates \( T(\varepsilon) \geq C \varepsilon^{-\frac{p-1}{\beta(p)}} \). This completes the proof of Theorem 1.
3 Proof of Theorem 2

In this section we are going to prove Theorem 2 by using a comparison argument for ordinary differential inequality of the second order (Kato’s lemma).

Let \( u \) be a local in time energy solution to (1) according to Definition 1 with lifespan \( T \) and let us fix \( t \in (0, T) \). If we choose a cut-off function \( \psi \in C_0^\infty([0, T) \times G) \) such that \( \psi = 1 \) on \([0, t] \times G \) in (3), then,

\[
\int_G \partial_t u(t, x) \, dx - \varepsilon \int_G u_1(x) \, dx = \int_0^t \int_G |u(s, x)|^p \, dx \, ds.
\]

Introducing the time-dependent functional

\[
U_0(t) = \int_G u(t, x) \, dx,
\]

we can rewrite the above equality in the following way

\[
U'_0(t) - U'_0(0) = \int_0^t \int_G |u(s, x)|^p \, dx \, ds.
\]

Then, \( U'_0 \) is differentiable with respect to \( t \) and

\[
U''_0(t) = \int_G |u(t, x)|^p \, dx \geq |U_0(t)|^p,
\]

where in the last step we applied Jensen’s inequality. We remark that

\[
U_0(0) = \varepsilon \int_G u_0(x) \, dx \geq 0 \quad \text{and} \quad U'_0(0) = \varepsilon \int_G u_1(x) \, dx \geq 0
\]

thanks to the assumptions on the initial data in the statement of Theorem 2. Then, employing Lemmas 2.1 and 2.2 in [17] (improved Kato’s lemma with upper bound estimate for the lifespan) to the functional \( U_0 \) we conclude the proof of Theorem 2.

4 Final remarks

Lately, the Cauchy problem for the semilinear damped wave equation with power nonlinearity has been studied in the framework of compact Lie groups in [9] (in this case the differential operator is the damped wave operator \( \partial_t^2 - L + \partial_t \)). In particular, in the compact case it has been proved for any exponent \( p > 1 \) the nonexistence of global in time solution, under certain sign assumptions for the initial data. This result is consistent with the one [14] for the semilinear heat equation on unimodular Lie group with polynomial volume growth. Indeed, we can read this result by saying that in the compact case the critical exponent for the semilinear damped wave equation is the Fujita exponent in the 0-dimensional case. Therefore, rather than the topological
dimension of the group $G$, it is the global dimension of $G$ (which is 0 for a compact Lie group) that provides the critical exponent. We refer to [2, Section II.4] for an overview on the growth properties of a Lie group.

A similar analysis can be done for the semilinear wave equation in (1). In the flat case, the subcritical condition $1 < p < p_{\text{Str}}(n)$ is equivalent to require that the quantity $\gamma(n, p) = -(n-1)p^2 + (n+1)p + 2$ is strictly positive. In the 0-dimensional case, we have $\gamma(0, p) > 0$ for any $p > 1$. Therefore, also for the semilinear Cauchy problem associated with wave operator $\partial_t^2 - \mathcal{L}$, it is possible to say that the global dimension of $G$ determines the range for $p$ in the blow-up result (in the sense that we have just explained).

Finally, we point out that not necessarily for all semilinear hyperbolic model on a compact Lie group a blow-up result can be proved for any $p > 1$. In the forthcoming paper [10], it will be shown that the combined presence of a damping term and of a mass term (0th order term with a positive multiplicative constant) modifies completely the situation since the global in time existence of small data solution can be proved in the evolution energy space without requiring any additional lower bound for $p > 1$.

References

1. Catania, D., Georgiev, V.: Blow-up for the semilinear wave equation in the Schwarzschild metric. Differ. Integral Equ. 19(7), 799–830 (2006)
2. Dungey, N., ter Elst, A.F.M., Robinson, D.W.: Analysis on Lie Groups with Polynomial Growth. Progress in Mathematics, vol. 214. Birkhäuser, Boston (2003)
3. Ebert, M.R., Reissig, M.: Regularity theory and global existence of small data solutions to semi-linear de Sitter models with power non-linearity. Nonlinear Anal. Real World Appl. 40, 14–54 (2018)
4. Fischer, V., Ruzhansky, M.: Quantization on Nilpotent Lie Groups, Progress in Mathematics, vol. 314. Springer, Birkhäuser (2016)
5. Galstian, A., Yagdjian, K.: Global in time existence of self-interacting scalar field in de Sitter spacetimes. Nonlinear Anal. Real World Appl. 34, 110–139 (2017)
6. Galstian, A., Yagdjian, K.: Finite lifespan of solutions of the semilinear wave equation in the Einstein-de Sitter spacetime. Rev. Math. Phys. (2020). https://doi.org/10.1142/S0129055X2050018X
7. Garetto, C., Ruzhansky, M.: Wave equation for sums of squares on compact Lie groups. J. Differ. Equ. 258(12), 4324–4347 (2015)
8. Lin, Y., Lai, N.-A., Ming, S.: Lifespan estimate for semilinear wave equation in Schwarzschild spacetime. Appl. Math. Lett. 99, 105997 (2020)
9. Palmieri, A.: On the blow—up of solutions to semilinear damped wave equations with power nonlinearity in compact Lie groups. J. Differ. Equ. 281, 85–104 (2021). https://doi.org/10.1016/j.jde.2021.02.002
10. Palmieri, A.: A global existence result for a semilinear wave equation with lower order terms on compact Lie groups. arXiv:2006.00759 (2020)
11. Palmieri, A.: Lifespan estimates for local solutions to the semilinear wave equation in Einstein–de Sitter spacetime. arXiv:2009.04388 (2020)
12. Palmieri, A.: Blow-up results for semilinear damped wave equations in Einstein–de Sitter spacetime. Z. Angew. Math. Phys. 72, 64 (2021). https://doi.org/10.1007/s00033-021-01494-x
13. Ruzhansky, M., Turunen, V.: Pseudo-differential operators and symmetries. In: Background Analysis and Advanced Topics, Volume 2 of Pseudo-Differential Operators. Theory and Applications. Birkhäuser, Basel (2010)
14. Ruzhansky, M., Yessirkegenov, N.: Existence and non-existence of global solutions for semilinear heat equations and inequalities on sub-Riemannian manifolds, and Fujita exponent on unimodular Lie groups. arXiv:1812.01933 (2018)
15. Ruzhansky, M., Yessirkegenov, N.: Hardy, Hardy–Sobolev, Hardy–Littlewood–Sobolev and Caffarelli–Kohn–Nirenberg inequalities on general Lie groups. arXiv:1810.08845v2 (2019)

16. Strauss, W.A.: Nonlinear scattering theory at low energy. J. Funct. Anal. 41(1), 110–133 (1981)

17. Takamura, H.: Improved Kato’s lemma on ordinary differential inequality and its application to semilinear wave equations. Nonlinear Anal. 125, 227–240 (2015)

18. Takamura, H., Wakasa, K.: The sharp upper bound of the lifespan of solutions to critical semilinear wave equations in high dimension. J. Differ. Equ. 251, 1157–1171 (2011)

19. Tsutaya, K., Wakaugi, Y.: Blow up of solutions of semilinear wave equations in Friedmann–Lemaître–Robertson–Walker spacetime. J. Math. Phys. 61, 091503 (2020). https://doi.org/10.1063/1.5139301

20. Yagdjian, K.: The semilinear Klein–Gordon equation in de Sitter spacetime. Discrete Contin. Dyn. Syst. Ser. S 2(3), 679–696 (2009)

21. Yagdjian, K.: Global existence of the scalar field in de Sitter spacetime. J. Math. Anal. Appl. 396(1), 323–344 (2012)

22. Yagdjian, K., Galstian, A.: Fundamental Solutions for the Klein–Gordon equation in de Sitter spacetime. Commun. Math. Phys. 285, 293–344 (2009)

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