Spiral-induced velocity and metallicity patterns in a cosmological zoom simulation of a Milky Way-sized galaxy

Robert J. J. Grand\textsuperscript{1,2}, Volker Springel\textsuperscript{1,2}, Daisuke Kawata\textsuperscript{3}, Ivan Minchev\textsuperscript{4}, Patricia Sánchez-Blázquez\textsuperscript{5,6}, Facundo A. Gómez\textsuperscript{7}, Federico Marinacci\textsuperscript{8}, Rüdiger Pakmor\textsuperscript{1} and David J. R. Campbell\textsuperscript{9}

\textsuperscript{1}Heidelberger Institut für Theoretische Studien, Schloss-Wolfsbrunnenweg 35, 69118 Heidelberg, Germany
\textsuperscript{2}Zentrum für Astronomie der Universität Heidelberg, Astronomisches Recheninstitut, Mönchhofstr. 12-14, 69120 Heidelberg, Germany
\textsuperscript{3}Mullard Space Science Laboratory, University College London, Holmbury St. Mary, Dorking, Surrey, RH5 6NT, United Kingdom
\textsuperscript{4}Leibniz-Institut für Astrophysik Potsdam (AIP), An der Sternwarte 16, D-14482, Potsdam, Germany
\textsuperscript{5}Departamento de Física Teórica, Universidad Autónoma de Madrid, Cantoblanco, E28049, Spain
\textsuperscript{6}Instituto de Astrofísica, Universidad Pontificia Católica de Chile, Av. Vicuña Mackenna 4860, Santiago, Chile
\textsuperscript{7}Max-Planck-Institut für Astrophysik, Karl-Schwarzschild-Str. 1, D-85748, Garching, Germany
\textsuperscript{8}Department of Physics, Kavli Institute for Astrophysics and Space Research, MIT, Cambridge, MA 02139, USA
\textsuperscript{9}Institute for Computational Cosmology, Department of Physics, Durham University, South Road, Durham, DH1 3LE, UK

Accepted XXX. Received YYY; in original form ZZZ.

ABSTRACT
We use a high resolution cosmological zoom simulation of a Milky Way-sized halo to study the observable features in velocity and metallicity space associated with the dynamical influence of spiral arms. For the first time, we demonstrate that spiral arms, that form in a disc in a fully cosmological environment with realistic galaxy formation physics, drive large-scale systematic streaming motions. In particular, on the trailing edge of the spiral arms the peculiar galacto-centric radial and azimuthal velocity field is directed radially outward and azimuthally forward, whereas it is radially inward and azimuthally backward at a given radius. Owing to the negative radial metallicity gradient, this systematic motion drives, at a given radius, an azimuthal variation in the residual metallicity that is characterised by a metal rich trailing edge and a metal poor leading edge. We show that these signatures are theoretically observable in external galaxies with Integral Field Unit instruments such as VLT/MUSE, and if detected, would provide evidence for large-scale systematic radial migration driven by spiral arms.

Key words: galaxies: evolution - galaxies: kinematics and dynamics - galaxies: spiral - galaxies: structure

1 INTRODUCTION
Spiral arms are typical features of late-type disc galaxies, which make up roughly $\sim 70\%$ of the bright galaxies in the local volume. They are found not only in the distribution of cold gas and young, bright stars, but also in the old stellar populations (Rix & Zaritsky 1995), which indicate that they are a dynamical phenomenon. Most theoretical models in the last 50 years have centred on variations of the classic density wave theory (Lin & Shu 1964) or swing amplification theory (Julian & Toomre 1966). In the former, spiral density enhancement is regarded as a crest of stars that preserves its shape as it propagates around the disc, whereas in the latter it is described by a shearing over-density that grows and decays around a preferred pitch angle (e.g., Baba et al. 2013; Grand et al. 2013; Michikoshi & Kokubo 2014). In the last decades many observational attempts have been made to test these theories, using methods such as the Tremaine-Weinberg equations (Tremaine & Weinberg 1984) and the spatial distribution of star forming tracers (e.g., Foyle et al. 2010; Ferreras et al. 2012). However, despite decades of study, the nature of spiral arms remains an unsolved problem in contemporary astrophysics.

In recent years, numerical simulations have provided new insights into the formation and evolution of spiral arms. Crucially, $N$-body simulations commonly show transient spi-
2 Numerical Simulation

We focus on one high-resolution cosmological zoom simulation taken from the Auriga simulation suite (see Grand et al. 2016, and Grand et al. in preparation for a full description), performed with the state-of-the-art magnetohydrodynamical moving-mesh code AREPO (Springel 2010).

The halo was initially selected from a parent dark matter simulation, which includes a wide range of galaxy formation physics. This implies that the nature of spiral arms is similar to the transient, winding spiral density enhancements commonly seen in idealised N-body simulations. Furthermore, we demonstrate for the first time that the systematic motion drives clear patterns of azimuthal variation of the metallicity distribution. We demonstrate that the peculiar velocity field and metallicity distribution around the spiral arms can be detected with current Integral Field Unit (IFU) instruments, such as VLT/MUSE (Bacon et al. 2010).

3 Results

Spiral structure is situated in the disc component of galaxies, and it follows that stars that spend much of their orbit in the disc mid-plane are dynamically responsive to such structures. We therefore focus our analysis on young stars (age < 3 Gyr) that belong to the thin disc, which constitute about 27% of the total stellar mass within a radius of 25 kpc. Note that if all star particles were to be considered, the observable features described in this Letter would be weaker, owing to contamination of bulge and halo stars. We focus on a single snapshot of galaxy Au 25 at a lookback time of 2.67 Gyr, which we choose because of its late time, quiescent environment and well-formed spiral arms in order to demonstrate clearly the dynamical signatures related to spiral structure, which is the aim of this study. The nature of the spiral arms and their evolution will be studied in a forthcoming paper.

3.1 Velocity fields

In the following, we define the azimuthal peculiar velocity, \( V_\phi \), as the difference between the azimuthal velocity of a star particle and the mean rotation velocity at the particle radius, and define \( V_\phi > 0 \) as faster than mean rotation. The radial and vertical peculiar velocities, \( V_R \) and \( V_Z \), are defined as the radial and vertical velocities, with \( V_R > 0 \) toward the galactic anti-centre and \( V_Z > 0 \) toward the north galactic pole.

In the left panel of Fig. 1 we show the face-on map of \( V_\phi \), with azimuthal over-density contours of the mass distribution, given by \( (\Sigma(R, \phi) - \Sigma(R))/\Sigma(R) \), overlaid in white contours. The spiral structure extends from about 5 to 15 kpc, and is accompanied by a well-defined spiral-shaped pattern in the azimuthal peculiar velocity field: stars rotate locally slower on the trailing side of the spiral arm, whereas they rotate locally faster on the leading side. The middle panel of Fig. 1 shows the face-on map of \( V_R \). Similarly, this velocity field reveals a spiral shaped pattern in which the spiral arm locus delineates the outward and inward streaming motions that are situated on the trailing and leading sides of the spirals, respectively. The right panel of Fig. 1 shows the face-on map of \( V_Z \), the fluctuations of which are of a lower amplitude in comparison to the planar velocity fields. We note that there may be indications of vertical modes in this galaxy (such as those shown in Gómez et al. 2016), which...
Figure 1. Face-on maps of the azimuthal (left), radial (middle) and vertical (right) peculiar velocity fields. Positive velocities are in the direction of rotation (azimuthal), the galactic anti-centre (radial) and positive vertical heights (vertical). Over-density contours of the mass distribution are indicated by the contours. The azimuthal peculiar velocity field is systematically slower (faster) on the trailing (leading) edge of the spiral, whereas the radial peculiar velocity points outward (inward) on the trailing (leading) edge. The amplitude of the fluctuations in the vertical peculiar velocity field are lower than the planar velocity fields, and show a less coherent pattern.

Figure 2. Face-on map (inclination, $i = 0$) of the azimuthal residual of the actual metallicity field (top panel) and that obtained for the case in which the metallicity field is artificially generated 120 Myr before the current time (bottom panel). Over-density contours of the mass distribution are indicated by the contours. At many radii there is an over-density of metal rich (poor) star particles on the trailing (leading) side of the spiral arm.

These patterns in the velocity fields are qualitatively similar to the systematic motions discussed in several recent studies of idealised simulations of isolated discs (e.g., Kawata et al. 2014; Hunt et al. 2015; Grand et al. 2015b) and in some observational work (e.g., Chemin et al. 2015). In the former studies, the systematic motions have been linked to transient, winding spiral density enhancements commonly seen in $N$-body simulations, in which star particles on the trailing or leading side of the spiral maintain their position with respect to the spiral peak, and are continuously torqued to larger or smaller guiding centre radii (Grand et al. 2012a, 2014).

3.2 Metal distribution

The torques applied to stars from spiral arms that generate the systematic streaming motions play a significant role in re-distributing individual stars around the disc, a process referred to as radial migration (Sellwood & Binney 2002; Minchev & Famaey 2010; Grand et al. 2012a). For a disc with a negative radial metallicity gradient, such as that of the Milky Way (e.g., Boeche et al. 2013; Bergemann et al. 2014; Anders et al. 2014), the radial re-distribution of stars can lead to changes in the metal distribution, in particular metal rich (poor) stars from the inner (outer) regions are brought to the outer (inner) regions. Such a radial re-distribution of stars has been shown to broaden the metallicity distribution (Minchev et al. 2013; Grand et al. 2015a), and is required to explain the large scatter in the age-metallicity relation in the solar neighbourhood (Chiappini et al. 2001; Haywood 2008; Casagrande et al. 2011). However, the observation of these trends is possible in the Milky Way only, for which star-by-star measurements are available.

A more directly observable signature of radial migration along spiral arms may come from azimuthal trends of metallicity at a given radius, in much the same spirit as the velocity trends are found in Fig. 1. To date, the only study of such a signature from a simulation perspective is that of Di Matteo et al. (2013), who studied the azimuthal variation of the metallicity distribution of old stars around a bar and found a pattern characteristic of radial migration...
around the bar co-rotation radius. A detailed study of the azimuthal variations of chemical abundances in APOGEE data, and their relation to Milky Way spiral structure, will be presented in Minchev et al. in preparation (see also Bovy et al. 2014).

In the top panel of Fig. 2, we show the face-on map of the azimuthal residual metallicity distribution, defined as \( \delta [\text{Fe/H}]/[R, \phi] = [\text{Fe/H}]/(R, \phi) - [\text{Fe/H}]/(R) \). It is clear that in the radial range of spiral structure the metallicity pattern is characterised by an over-density of metal rich stars on the trailing side of the spiral arm, whereas an over-density of metal poor stars are found on the leading side of the spiral. These features are a consequence of the radial metallicity distribution, which has a radial gradient of \(-0.035 \text{ dex kpc}^{-1}\) and metallicity dispersion of about 0.19 dex at a given radius. We note that the trend is even more clear in the bottom panel of Fig. 2, in which the radial metallicity distribution is artificially set 120 Myr earlier (about a dynamical time) with a fixed radial gradient of \(-0.08 \text{ dex kpc}^{-1}\) and a radially constant metallicity dispersion of 0.05 dex. This confirms that the dynamics are consistent with large-scale radial migration along the spiral arms, and also that the signatures become more clear for steeper radial gradients and narrower metallicity dispersions, respectively. This is the first time that such a trend has been shown, and provides a further observational test for the radial migration driven by spiral arms.

### 3.3 Observing the line-of-sight signatures

The features of the peculiar velocity and residual metallicity fields shown above can be directly observed in external galaxies with IFU instruments, such as VLT/MUSE. To demonstrate how the peculiar velocity field and metallicity distribution are mapped to the line-of-sight (LOS) velocity field\(^1\), \(\delta V_{\text{los,pec}}\), and LOS metallicity distribution, we set the disc to an inclination of 30 degrees, which is enough inclined to observe planar peculiar velocities while maintaining a clear view of the spiral structure. The projected LOS peculiar velocity map and residual metallicity map are shown in the top panels of Fig. 3. For \(X < 0\), positive \(V_{\phi}\) is mapped to positive LOS velocities (away from the observer), and to negative LOS velocities (toward the observer) for \(X > 0\). Positive \(V_R\) is mapped to positive LOS velocities for \(Y > 0\), and negative LOS velocities for \(Y < 0\).

To quantify the fluctuations in both fields, we show in the bottom panels of Fig. 3 the azimuthal profiles of the residuals of the mass surface density, metallicity and LOS peculiar velocity fields at two different radii. We define \(\phi = 0\) for \(X = 0\) and \(Y > 0\), which increases counter-clockwise in the disc plane. The patterns in the LOS peculiar velocity, \(\delta V_{\text{los,pec}}\), depend on the location of the spiral arms, because the direction of the LOS projection of the radial and azimuthal peculiar velocities flips around \(Y = 0\) and \(X = 0\), respectively. At \(R = 10\) kpc, the spiral arm at \(\phi \sim 5.5\) rad is located at \(Y > 0\). Both the leading (\(\phi < 5.5\) rad) and trailing side (\(\phi > 5.5\) rad), have positive \(X\) and \(Y\) coordinates, where azimuthally fast (slow) and radially inward (outward) peculiar velocities on the leading (trailing) edge give rise to negative (positive) \(\delta V_{\text{los,pec}}\). Clockwise (decreasing \(\phi\)) of this spiral arm the \(\delta V_{\text{los,pec}}\) remains negative because the radially positive motion on the trailing edge of the next spiral (\(\phi \sim 3.0\) rad) for \(Y < 0\) is directed toward the observer. On the leading edge of this spiral arm (\(\phi \sim 1.5\) rad) the large peculiar azimuthal velocity leads to a positive \(\delta V_{\text{los,pec}}\). A similar trend is present at \(R = 16\) kpc, though there is a small additional peak at \(\phi \sim 4.0\) owing to the phase shift in spiral arm position with respect to \(R = 10\) kpc that causes negative peculiar azimuthal velocities on the trailing side of the spiral to contribute to positive \(\delta V_{\text{los,pec}}\) for \(X > 0\) and \(Y < 0\). The semiamplitude of \(\delta V_{\text{los,pec}}\) is about \(10 - 15 \text{ km s}^{-1}\). For the metallicity residuals, the semiamplitude is \(\sim 0.05\) dex, which yields a total variation of \(\sim 0.1\) dex. The semiamplitudes of the LOS peculiar velocity can be increased up to \(\sim 20 \text{ km s}^{-1}\) for inclinations up to \(60\) degrees, however the spatial resolution is lower. The magnitudes of these fluctuations should be large enough to be detected with IFU observations of nearby late-type galaxies.

### 4 CONCLUSIONS

In this study we have analysed signatures in the peculiar velocity and residual metallicity fields linked to the dynamical influence of spiral arms, in one of the high resolution, fully cosmological zoom simulations from the Auriga suite

---

\(^1\) We note that the contribution from the vertical peculiar velocity to the LOS velocity within this radial range is minor and does not affect the trends discussed in this paper.
We have demonstrated that the peculiar azimuthal velocity is locally slower (faster) on the trailing (leading) edge of the spiral arm. Similarly, the peculiar radial velocity is directed radially outward (inward) on the trailing (leading) edge of the spiral arm, which in combination with the azimuthal velocity creates a systematic streaming motion along the spiral arm. This represents the first confirmation of systematic radial migration around spiral arms in a fully cosmological zoom simulation, which is the first confirmation of systematic radial migration around apeculiar streaming motion along the spiral arm. This represents a combination with the azimuthal velocity creates a system-}

In addition, we show for the first time that the radial migration caused by the spiral arms leads to azimuthal variations of the metallicity distribution: at a given radius, star particles that originated from inferior regions of the disc are metal rich with respect to the azimuthal mean metallicity at that radius, because of the negative metallicity gradient commonly observed in disc galaxies. As indicated by the systematic streaming motions, the metal rich particles are transported outward along the trailing edge of the spiral. Similarly, the metal poor star particles that originate in the outer disc regions are transported radially inwards along the leading edge of the spiral. The result of these motions is a residual metallicity pattern in azimuth which is systematically more metal rich (poor) along the trailing (leading) edge of the spiral arm at many radii.

Finally, we have shown that the azimuthal variations of the peculiar LOS velocity and metallicity maps in a disc inclined at an angle of 30 degrees are about 20-30 km s$^{-1}$ and 0.1 dex, respectively. These variations can be detectable in nearby late-type spiral galaxies with IFU instruments such as VLT/MUSE.

We note that a systematic difference in azimuthal velocity across a density wave-like spiral arm has been suggested by Minchev & Quillen (2008); Pasetto et al. (2015). However, these motions would depend on the location of the resonance points and are not expected to drive a metallic-}

This Letter is an encouraging first step toward mak-}

search Council under ERC-StG grant EXAGAL- 308037. Part of the simulations of this paper used the SuperMUC system at the Leibniz Computing Centre, Garching, under the project PR85JE of the Gauss Centre for Supercomputing. This work used the DiRAC Data Centric system at Durham University, operated by the Institute for Computational Cosmology on behalf of the STFC DiRAC HPC Facility ‘www.dirac.ac.uk’. This equipment was funded by BIS National E-infrastructure capital grant ST/K00042X/1, STFC capital grant ST/H008519/1, and STFC DiRAC Operations grant ST/K003267/1 and Durham University: DiRAC is part of the National E-Infrastructure.

REFERENCES
Anders F., et al., 2014, A&A, 564, A115
Baba J., Saitoh T. R., Wada K., 2013, ApJ, 763, 46
Bacon R., et al., 2010, in Society of Photo-Optical In-
strumentation Engineers (SPIE) Conference Series. p. 8, doi:10.1117/12.856027
Bergemann M., et al., 2014, A&A, 565, A89
Boeche C., et al., 2013, A&A, 559, A59
Bovy J., et al., 2014, ApJ, 790, 127
Casagrande L., Schönrich R., Asplund M., Cassisi S., Ramírez I., Meléndez J., Bensby T., Feltzing S., 2011, A&A, 530, A138
Chemin L., Renaud F., Soubiran C., 2015, A&A, 578, A14
Chiappini C., Matteucci F., Romano D., 2001, ApJ, 554, 1044
Di Matteo P., Haywood M., Combes F., Semelin B., Naughton P. A., 2013, A&A, 553, A102
Ferreras I., Cropper M., Kawata D., Page M., Huzenstien W. A., 2012, MNRAS, 424, 1636
Foyle K., Rix H.-W., Walter F., Leroy A. K., 2010, ApJ, 725, 534
Gómez F. A., White S. D. M., Marinacci F., Slater C. T., Grand R. J. J., Springel V., Pakmor R., 2016, MNRAS, 456, 2779
Grand R. J. J., Kawata D., Cropper M., 2012a, MNRAS, 421, 1529
Grand R. J. J., Kawata D., Cropper M., 2012b, MNRAS, 426, 167
Grand R. J. J., Kawata D., Cropper M., 2013, A&A, 553, A77
Grand R. J. J., Kawata D., Cropper M., 2014, MNRAS, 439, 623
Grand R. J. J., Kawata D., Cropper M., 2015a, MNRAS, 447, 4018
Grand R. J. J., Bovy J., Kawata D., Hunt J. A. S., Famaey B., Siebert A., Monari G., Cropper M., 2015b, MNRAS, 453, 1867
Grand R. J. J., Springel V., Gómez F. A., Marinacci F., Pakmor R., Campbell D. J. R., Jenkins A., 2016, MNRAS, 368, 1175
Haywood M., 2008, MNRAS, 388, 1715
Hunt J. A. S., Kawata D., Grand R. J. J., Minchev I., Pasetto S., Cropper M., 2015, MNRAS, 450, 2132
Julian W. H., Toomre A., 1966, ApJ, 146, 810
Kawata D., Hunt J. A. S., Grand R. J. J., Pasetto S., Cropper M., 2014, MNRAS, 443, 2757
Lin C. C., Shu F. H., 1964, ApJ, 140, 646
Marinacci F., Pakmor R., Springel V., 2014, MNRAS, 437, 1750
Michikoshi S., Kokubo E., 2014, ApJ, 787, 174
Minchev I., Famaey B., 2010, ApJ, 722, 112
Minchev I., Quillen A. C., 2008, MNRAS, 386, 1579
Minchev I., Chiappini C., Martig M., 2013, A&A, 558, A9
Monari G., Famaey B., Siebert A., 2016, MNRAS, 457, 2569
Pasetto S., Natale G., Kawata D., Chiosi C., Hunt J. A. S., 2015, preprint, (arXiv:1512.05367)
Planck Collaboration et al., 2014, A&A, 571, A16
Rix H.-W., Zaritsky D., 1995, ApJ, 447, 82
Sellwood J. A., 2011, MNRAS, 410, 1637
Sellwood J. A., Binney J. J., 2002, MNRAS, 336, 785
Springel V., 2010, MNRAS, 401, 791
Grand et al.

Tremaine S., Weinberg M. D., 1984, ApJ, 282, L5
Vogelsberger M., Genel S., Sijacki D., Torrey P., Springel V.,
   Hernquist L., 2013, MNRAS, 436, 3031