Characterization and Optimization of Coherent MZI-Based Nanophotonic Neural Networks Under Fabrication Non-Uniformity

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Abstract—Silicon-photonic neural networks (SPNNs) and artificial intelligence (AI) accelerators have emerged as promising successors to electronic accelerators by offering orders of magnitude lower latency and higher energy efficiency. Nevertheless, the underlying silicon photonic devices in SPNNs are sensitive to inevitable fabrication-process variations (FPVs) stemming from optical lithography imperfections. Consequently, the inferencing accuracy in an SPNN can be highly impacted by FPVs—e.g., can drop to below 10%—the impact of which is yet to be fully studied. In this article, we, for the first time, model and explore the impact of optical phase noise due to FPVs in the waveguide width, silicon-on-insulator (SOI) thickness, and etch depth in coherent SPNNs that use Mach–Zehnder interferometers (MZIs). Leveraging such models, we propose a novel variation-aware, design-time optimization solution to improve MZI tolerance to different FPVs in SPNNs. Simulation results for two example SPNNs of different scales under realistic and correlated FPVs indicate that using the optimized MZIs can lead to significant improvements in the network inferencing accuracy. The proposed one-time optimization imposes low area overhead and hence is applicable even to resource-constrained designs.

Index Terms—Silicon photonic integrated circuits, fabrication-process variations, deep learning, optical neural networks.

I. INTRODUCTION

A Machne learning algorithms are being utilized in a wide range of applications ranging from autonomous driving, real-time speech translation, and network anomaly detection to pandemic growth and trend prediction. With the rising demand for advanced neural networks to address even more complex problems, artificial intelligence (AI) accelerators need to consistently deliver better performance and improved accuracy while being energy-efficient. Unfortunately, in the post-Moore’s law era, electronic AI accelerator architectures face fundamental limits in their processing capabilities due to the limited bandwidth and low energy efficiency of metallic interconnects. Silicon photonics can alleviate these bottlenecks by enabling chip-scale optical interconnects with ultra-high bandwidth, low-latency, and energy-efficient communication, and light-speed chip-scale optical computation [1]. Leveraging silicon-photonic-enabled optical interconnect and computation, many integrated silicon-photonic neural networks (SPNNs) have been recently proposed [1].

Prior work has shown that the complexity of matrix-vector multiplication can be reduced from $O(N^2)$ to $O(1)$ in SPNNs [2]. Also, as SPNNs use photons for computation, there is negligible latency associated with inferencing. Such benefits have positioned SPNNs as a promising alternative to traditional electronic neural networks [3]. Nevertheless, fabrication-process variations (FPVs) due to inevitable optical lithography imperfections lead to undesired changes in the critical dimensions of SPNNs’ building blocks (e.g., waveguide width and thickness in Mach–Zehnder interferometers (MZIs) in coherent SPNNs), imposing inaccuracies during matrix-vector multiplication. In [4], we have shown that random uncertainties due to FPVs and thermal crosstalk can result in up to a catastrophic 70% accuracy loss, even in mature fabrication processes. Existing approaches for improving the resilience of SPNNs against FPVs largely rely on compensating their impact by individually calibrating MZIs in the network [5], [6]. However, such solutions impose additional calibration (i.e., tuning) power overhead and are not scalable as their complexity grows with the number of devices in SPNNs.

The main contribution of this article is in developing, to the best of our knowledge, the first comprehensive analysis of the impact of physical-level FPVs on coherent SPNN performance. We consider variations in the waveguide width, silicon-on-insulator (SOI) thickness, and etch depth, which impacts the slab thickness, based on realistic FPV maps developed using different correlation lengths in the variations and experimental
measurements. We model the impact of FPVs at the device level for MZI performance and at the network level for arrays of cascaded MZIs in SPNNs. At the system level, we explore the impact of variations on SPNN inferencing accuracy while considering different FPVs and variation correlation lengths. Leveraging our detailed device-level models, we also develop a novel design optimization solution to improve MZI performance in SPNNs under different FPVs. Our simulation results for two example SPNNs (with 1380 and 20,580 phase shifters) under FPVs show that while inferencing accuracy can drop to as low as 7.73% under different variations, using our optimized MZIs can significantly improve the inferencing accuracy (e.g., by up to 72% on average in a large SPNN). Note that this article does not consider the impact of thermal crosstalk and FPVs in directional couplers in MZIs.

The rest of the article is organised as follows. Section II presents a background on MZIs and SPNNs and a summary of prior related work. In Section III, we analyze the impact of FPVs on MZIs (device level) and array of cascaded MZIs (network level) in SPNNs. Section IV presents our MZI design optimization solution to design devices with improved tolerance to different FPVs. Section V presents simulation results that highlight how the optimized MZIs can improve the accuracy of the unitary transformation (at the layer level) and the inferencing accuracy (at the system level), in the presence of FPVs. Last, we draw conclusions in Section VI.

II. BACKGROUND AND PRIOR RELATED WORK

This section summarizes fundamentals of MZIs and coherent SPNNs designed based on MZIs. Also, it presents some preliminary models to help understand the impact of FPVs on photonic waveguides, and reviews some prior related work.

A. Mach–Zehnder Interferometer (MZI)

A $2 \times 2$ MZI in an SPNN consists of two tunable phase shifters ($\phi$ and $\theta$) on the upper arm and two 50:50 beam splitters as shown in Fig. 1(c). The phase shifters are used to apply configurable phase shifts and obtain varying degrees of interference between the optical signals traversing the MZI arms. In SPNNs, phase shifters are often implemented using thermal micro-heaters for lower optical loss, where the effective index of the underlying waveguide changes with temperature (i.e., due to thermo-optic effect of silicon), hence altering the phase of the optical signal. Note that the proposed analyses in this article are independent of the phase-shift mechanism in the MZI. Moreover, $2 \times 2$ beam splitters can be designed using directional couplers (DCs) where a fraction of the optical signal (defined as $\kappa$) at an input port is transmitted to an output port, and the remaining signal is coupled to the other output port, as shown in the inset of Fig. 1(c). The ideal transfer matrix ($T_{MZI}$) for an MZI with two phase shifters ($\theta$ and $\phi$) and two 50:50 beam splitters can be defined as [7] (Fig. 1(c)):

$$T_{MZI}(\theta, \phi) = \begin{pmatrix}
\frac{1}{2} (e^{i\theta} - 1) & \frac{i}{2} (e^{-i\theta} + 1) \\
\frac{i}{2} (e^{i\theta} + 1) & \frac{1}{2} (e^{-i\theta} - 1)
\end{pmatrix}.$$  (1)

Such a fraction of the optical signal (defined as $\kappa$) can be mathematically modeled as $\kappa = e^{i\phi} = \frac{1}{2} (e^{-i\theta} + 1)$ where $\phi$ and $\theta$ are the unitary matrices with dimension of $m \times m$. A unitary matrix can be realized by using a cascaded array of $2 \times 2$ MZIs. As shown in Fig. 1(a) and (b), this unit is called the optical-interference unit (OIU) and is responsible for performing the MAC operation.

Several approaches have been proposed to design the architecture of MZI arrays to perform MAC operations in the optical domain [8], [9], [10]. Out of these, the Clements design, due its symmetric nature, has a low optical loss and footprint. Therefore,
we use the Clements design [9] to transform each unitary matrix \( \left( U_{m,n}^{T} \times m \right) \) and \( V_{m,n}^{H} \times m \) to a cascaded MZI array with a specific phase setting per MZI in the network. In the Clements method, each unitary matrix will be mapped to a cascaded MZI array with a total number of \( \frac{N(N-1)}{2} \) MZIs, where \( N \) is the size of the designated unitary matrix. Note that the diagonal matrix \( \left( \Lambda_{m,n}^{T} \times m \right) \) can be realized with an array of MZIs with one input and one output terminated. The optical-amplification unit (OAU) in Fig. 1(a), which is required to obtain arbitrary diagonal matrix, can be realized using semiconductor optical amplifiers (SOAs).

C. Fabrication-Process Variations (FPVs)

FPVs in silicon photonics originate in optical-lithography process imperfections, contributing to changes in the waveguide width, SOI thickness (dominated by the host wafer), and slab thickness (in case of a ridge waveguide). Such changes deviate the effective index \( n_{eff} \) in a waveguide and in turn the propagation constant \( \beta \) which determines the optical phase of the signal traversing the waveguide. Considering Fig. 1(c) and as an example, the effective index \( n_{eff} \) in a ridge waveguide depends on the optical wavelength and the critical dimensions of the waveguide [11], i.e., width \( w \), SOI thickness \( t \), and slab thickness \( h \) in Fig. 1(c). Note that \( h = 0 \) for a strip waveguide. The relation can be defined as:

\[
\beta(\lambda, w, t, h) = \left( \frac{\lambda}{2\pi} \right) \beta(\lambda, w, t, h),
\]

where \( \lambda \) is the optical wavelength. Here, \( \beta \) can be defined as \( \beta = \psi/L \), where \( \psi \) is the single pass phase-shift induced in a waveguide of length \( L \). Leveraging (2), propagation constant changes \( \Delta \beta \) in a ridge waveguide under FPVs is given by:

\[
\Delta \beta = 2\pi \left( \frac{\partial n_{eff}}{\partial w} \rho_{w} + \frac{\partial n_{eff}}{\partial t} \rho_{t} + \frac{\partial n_{eff}}{\partial h} \rho_{h} \right).
\]

Here, \( \rho_{w}, \rho_{t}, \) and \( \rho_{h} \) are the variations in the waveguide width, SOI thickness, and slab thickness, respectively. When using a strip waveguide, \( \frac{\partial n_{eff}}{\partial t} = \rho_{h} = 0 \). Note that, in this article, we do not consider the variations in \( L \) (see Section IV).

D. Related Work on FPV Analysis in SPNNs

Our prior work in [4] studied the impact of random phase noise due to FPVs and thermal crosstalk at the system level in SPNNs by developing a framework that identifies critical components in the network. In [7], imprecisions were introduced in SPNNs after software training such that pre-fabrication training tends to counter the impact of both FPVs and thermal crosstalk at the system level in SPNNs.

III. MODELING FPVs IN COHERENT SPNNs

This section presents a detailed bottom-up analysis of the impact of FPVs in the waveguide width, SOI thickness, and slab thickness at the device level (i.e., MZI devices) and network level (OIU in Fig. 1(b)) in coherent SPNNs. We show the impact of FPVs on SPNN inferencing accuracy (system level) in Section V.

A. Device Level: MZI Performance Under FPVs

To study the impact of FPVs on MZI devices, we should first model FPVs in silicon photonics and explore how MZI devices experience such FPVs. In our prior work [13], we have developed realistic wafer variation maps that model radial-variation effects and correlation among different variations—both of which are critical to realistically model FPVs in silicon photonics—in SOI wafers. Such maps were developed based on mean, standard deviation (\( \sigma \)), and correlation lengths (\( l \)) experimentally characterized in collaboration with CEA-Leti in [13]. Table I summarizes different parameters considered to generate FPV maps for our calculations, which are based on analyzing experimental data from characterizing actual 200-mm wafers at CEA-Leti. Fig. 2-top shows examples of wafer and die maps generated using our in-house FPV wafer map simulator with a resolution of \( 10 \mu m \times 10 \mu m \) (i.e., the mesh size in the map—see Fig. 2). MZIs are bulky devices (e.g., \( \approx 340 \mu m \) in length [14]), and hence every section of the device will experience slightly different FPVs (see Fig. 2). Such a difference requires characterization of each phase shifter and passive splitter in the photonic circuit while relying on detectors in the output to calibrate the parameters.

The aforementioned methods focus on mitigating deviations in SPNNs post-fabrication by either post-fabrication training methods to compensate for any additionally introduced phase noises, which might impact the network accuracy [5], or calibrating each noisy component, where additional error-detection phases are required and the complexity can increase as the SPNN scales up [6]. In this work, we focus on exploring and optimizing the physical design of MZI devices in coherent SPNNs under FPVs prior to fabrication, hence improving the network tolerance under different FPVs. We show that by exploring and optimizing MZI device physical-level design, we can improve the relative-variation distance (RV D) in SPNNs, which quantifies the deviation between the intended unitary matrix and the deviated unitary matrix, to enhance the overall inferencing accuracy.
devices. Considering the MZI shown in Fig. 2 with its four arms labeled (Arm1–Arm4), we average the width, SOI thickness, and slab thickness (when using a ridge waveguide in the MZI) variations observed on each MZI arm, separately over the section colored on each arm.

Considering the MZI in Fig. 2, the optical signals traversing the two opposite arms of the device—before the input DC (i.e., Arms 1 and 2) and output DC (i.e., Arms 3 and 4)—should only experience the desired phase change \( \phi \) or \( \theta \), adjusted after the network training. Note that the phase shifters are integrated on the top of the silicon waveguides. However, due to the impact of non-uniform FPVs on each individual arm, the optical signals experience some undesired phase changes. Assuming that each arm’s length \( L_1 = L_2 = L_\theta \) and \( L_3 = L_4 = L_\theta \) does not undergo any variations, the optical phase difference between the two optical signals traversing the opposite arms (Arms 1–2 and Arms 3–4) and interfering at the input of DC1 (\( \Delta \Phi_{DC1} \)) and DC2 (\( \Delta \Phi_{DC2} \)) is:

\[
\Delta \Phi_{DC1} = \phi + |\Delta \beta_1 L_1 - \Delta \beta_2 L_2|, \\
\Delta \Phi_{DC2} = \theta + |\Delta \beta_3 L_3 - \Delta \beta_4 L_4|.
\] (4a)

Here, \( \Delta \beta_1, \Delta \beta_2, \Delta \beta_3, \) and \( \Delta \beta_4 \) are, respectively, the propagation constant changes on MZI’s Arms 1–4, which can be calculated using (3).

Leveraging (3) and (4), and assuming variation-free DCs, the MZI transfer matrix in (1) can be updated to take into consideration the impact of FPVs on MZI arms, resulting in undesired optical phase noises:

\[
T_{MZI}(\theta, \phi) = \begin{pmatrix}
\frac{\sqrt{2}}{2} e^{i(\theta + \Delta \beta_1 L_1)} & \frac{\sqrt{2}}{2} e^{i(\Delta \beta_1 L_4)} \\
\frac{\sqrt{2}}{2} e^{i(\theta + \Delta \beta_2 L_2)} & \frac{\sqrt{2}}{2} e^{i(\Delta \beta_2 L_4)}
\end{pmatrix} \\
\begin{pmatrix}
\frac{\sqrt{2}}{2} e^{i(\phi + \Delta \beta_3 L_1)} & \frac{\sqrt{2}}{2} e^{i(\Delta \beta_3 L_2)} \\
\frac{\sqrt{2}}{2} e^{i(\phi + \Delta \beta_4 L_3)} & \frac{\sqrt{2}}{2} e^{i(\Delta \beta_4 L_2)}
\end{pmatrix},
\] (5)

where, as shown in Fig. 2, \( L_1 = L_2 = L_\theta \) and \( L_3 = L_4 = L_\theta \). In this article, we assume \( L_\theta = L_\phi \approx 135 \, \mu m \), considered as an example based on [14]. Leveraging (5), we can capture the impact of non-uniform variations in the waveguide width, SOI thickness, and slab thickness in any MZI design in coherent SPNNs. Note that FPVs will deviate the splitting ratio, which ideally should be 50:50, in the input and output directional couplers (DC1 and DC2 in Fig. 2), hence affecting the network accuracy [4]. However, this article focuses on the optical phase noise due to FPVs in MZIs and considers variation-free DCs in SPNNs. Also, note that the analyses proposed in this section are independent of the example MZI considered in Fig. 2.

### B. Network Level: OIU Performance Under FPVs

Here, we model the impact of FPVs on the performance of an OIU, shown in Fig. 1(b). Note that in this article we do not consider the impact of FPVs in the OAUs and NAUs (see Fig. 1(a)) as they are often implemented either electronically or opto-electronically [2], [15]. Recall from Section II-B that, given a weight matrix \( W = U \Sigma V^H \), the matrices \( U, \Sigma, \) and \( V \) can be decomposed to \( U \Sigma V \) and \( \Sigma V \) phases on each MZI in an OIU using Clements decomposition [9]. Under FPVs, the transfer matrix of each MZI in the OIU will deviate in a manner that can be calculated using (5). To analyze the impact of such variations at the network level, we use relative-variation distance (\( RVD \)). \( RVD \) determines the deviation between an intended matrix and a deviated matrix [4]. We found that there is a strong correlation between network-level \( RVD \) and system-level inferencing accuracy in SPNNs (we will discuss this in Section V). \( RVD \) can be defined as:

\[
RVD(W, \overline{W}) = \frac{\sum_m \sum_n |W_{m,n} - \overline{W}_{m,n}|}{\sum_m \sum_n |W_{m,n}|},
\] (6)

where \( |.| \) denotes the absolute value of a complex number. \( W \) is the nominal weight matrix and \( \overline{W} \) is the deviated weight matrix under FPVs related to a fully connected layer. \( W_{m,n} \) denotes the element at the \( m^{th} \) row and \( n^{th} \) column of \( W \). Each MZI in the OIU has a unique impact on each element of the weight matrix. Accordingly, variations related to each MZI in the network have a unique effect on the overall \( RVD \). Higher \( RVD \) means the actual weight matrix is more deviated from the intended one, which can be interpreted as observing a lower inferencing accuracy at the system level.

To compute \( \overline{W} \) in (6), we first need to analyze the impact of non-uniform FPVs in the waveguide width, SOI thickness, and slab thickness on each individual MZI in an OIU. As the first step, we calculate the total dimension of an OIU based on the length of an individual MZI \( l_{MZI} \) and the distance between its input and output ports \( g_{MZI} \). In this work and as an example, we consider \( l_{MZI} = 340 \, \mu m \) and \( g_{MZI} = 30 \, \mu m \), and based on the phase shifter length in [14] (i.e., \( \approx 135 \, \mu m \)). Accordingly and using the Clements design for the OIU [9], the total dimension of the OIU can be calculated. As the second step, we use our in-house FPV wafer map simulator (see Section III-A) to generate a die map that matches the size of OIU. We then place the OIU on the die map and extract FPV information for each individual MZI in the OIU. Note that each MZI itself experiences different variations (non-uniform variations across a single device; a.k.a. intra-device variations), and the FPVs between two different MZIs are also different. All such non-uniformities are considered in our device-level (Section III-A) and network-level analyses. By capturing FPVs for each individual MZI in an OIU,
we can calculate the deviated $W$ in (6) based on:
\[
\bar{W}_m = \bar{U}_m \times \Sigma_m \times \bar{V}_m^H.
\]
Here, $\bar{U}_m$ and $\bar{V}_m^H$ are the deviated unitary transfer matrices and $\Sigma_m$ is the deviated diagonal matrix under FPVs. They can be calculated based on multiplying MZI transfer matrices under FPVs (the model in (5)), and in a specific order determined by the Clements design [9]. For example, for $\bar{U}_m$, we have:
\[
\bar{U}_m = D \left( \prod_{(m,n) \in S} T_{MZI,m,n}' \right),
\]
where $m$ and $n$ should be calculated based on the mapping method (e.g., Clements [9]) used to map the weight matrices to cascaded MZI arrays. Also, $S$ is the order of multiplication which again should be determined by the mapping method. Moreover, $D$ is a diagonal matrix with unity magnitude and is not related to the physical placement of MZIs. Similarly, we can calculate $\bar{V}_m^H$ and $\Sigma_m$. Although we considered the Clements method for mapping the weights to phase settings of a cascaded MZI array in OIU's, our network-level models in this section can work with any mapping method.

IV. SPNN DESIGN OPTIMIZATION UNDER FPVs

In this section, we explore the design space of MZI devices under different FPVs to optimize their performance in SPNNs. In particular, we focus on minimizing the impact of FPVs on MZI arms that imposes undesired optical phase noises in the device, leading to faulty matrix-vector multiplication. As discussed in Section III, FPVs also deviate the splitting ratio in DCs in an MZI. Nevertheless, the design optimization solution in this section assumes ideal DCs. Note that FPV-tolerant DCs can be designed based on the method proposed in [16].

Considering (4), one can alleviate the impact of FPV-induced optical phase noise in an MZI by implying $|\Delta \beta_1 L_1 - \Delta \beta_2 L_2| \to 0$ and $|\Delta \beta_3 L_3 - \Delta \beta_4 L_4| \to 0$. Accordingly, to obtain a phase-noise-free MZI, we should have:
\[
\frac{L_1}{L_2} = \frac{\Delta \beta_2}{\Delta \beta_1} \quad \text{and} \quad \frac{L_3}{L_4} = \frac{\Delta \beta_4}{\Delta \beta_3}.
\]
This indicates that the length ratio between any two opposite arms in the MZI should be inversely proportional to the ratio of the changes in their waveguide propagation constants ($\Delta \beta$), under non-uniform FPVs. In this section and for brevity, we assume $L_1 = L_2 = L_3 = L_4$ and without variations. As a result and based on (9), we should minimize $|\Delta \beta_1 - \Delta \beta_2|$ and $|\Delta \beta_3 - \Delta \beta_4|$ under different FPVs. In other words, we should make sure that the propagation constant changes on the two opposite arms in an MZI are as small as possible (i.e., $\Delta \beta_1 \to 0$, $\Delta \beta_2 \to 0$, $\Delta \beta_3 \to 0$, and $\Delta \beta_4 \to 0$), or the propagation constant changes on the two opposite arms are as close as possible (i.e., $\Delta \beta_1 \to \Delta \beta_2$ and $\Delta \beta_3 \to \Delta \beta_4$).

Considering (3), $\Delta \beta$ is proportional to the rate of changes in the waveguide’s effective index under FPVs (i.e., $\frac{\partial n_{eff}}{\partial \lambda}$), where $X$ denotes the design parameter under FPVs; i.e., $X = w$ for width, $X = t$ for SOI thickness, and $X = h$ for slab thickness variations. In our prior work [13], we found that as the waveguide width increases, $\frac{\partial n_{eff}}{\partial w}$ decreases, especially under waveguide width variations (i.e., when $X = w$). This is because the waveguide width increases, a bigger portion of the optical mode is confined in the waveguide core (and the confinement is also stronger), and hence the variations in the waveguide width will create less distortion in the optical mode in the waveguide. Also, note that increasing the waveguide width helps reduce the propagation loss in strip and ridge waveguides [13]. Fig. 3(a) shows the rate of changes in the effective index of a strip waveguide with $t = 220$ nm (see Fig. 2) and when the waveguide width ($w$) changes from 350 nm to 1200 nm, both considered as an example. As it can be seen, as the waveguide width increases, $\frac{\partial n_{eff}}{\partial w}$ decreases sharply but $\frac{\partial n_{eff}}{\partial t}$ decreases slightly and stays higher than $\frac{\partial n_{eff}}{\partial w}$ under different waveguide widths. While waveguide width can be changed during the design time, the SOI thickness cannot be changed; this parameter is determined by the host SOI wafer.

As it can be seen from Fig. 3(a), both $\frac{\partial n_{eff}}{\partial h}$ and $\frac{\partial n_{eff}}{\partial w}$ are still high in strip waveguides. To address this, we also explore the design of MZIs using a shallow-etched ridge waveguide with a slab thickness ($h$) of 150 nm, as shown in Fig. 2. We simulated different slab thicknesses from 60 nm to 180 nm (results are not shown in the article), and $h = 150$ nm returned the best results. Such a ridge waveguide is common in the design of grating structures [17]. By adding the slab to a strip waveguide (i.e., making it a ridge waveguide), the optical mode is pulled mostly towards the slab region, hence SOI thickness variations should have less impact on the optical mode. Similar to Fig. 3(a), Fig. 3(b) shows the rate of changes in the effective index of the shallow-etched ridge waveguide with $t = 220$ nm and $h = 150$ nm, and when the waveguide width ($w$) changes from 350 nm to 1200 nm. Observe that, compared to the strip waveguide, both $\frac{\partial n_{eff}}{\partial h}$ and $\frac{\partial n_{eff}}{\partial w}$ are smaller in the ridge waveguide. The shallow-etched ridge waveguide also suffers from variations in its slab thickness. Nevertheless, as it is shown by Fig. 3(b), $\frac{\partial n_{eff}}{\partial h}$ is much smaller than $\frac{\partial n_{eff}}{\partial w}$ in the shallow-etched ridge waveguide, and it decreases sharply as the waveguide width increases.

Considering the results in Fig. 3, designing MZIs with wider strip and shallow-etched ridge waveguides should help minimize
the changes in the propagating constant on each MZI arm (see (3) and (9)). Moreover, to increase the waveguide width, an important design consideration is to include waveguide tapers on the MZI arms as shown in Fig. 2. Waveguide tapers are essential to avoid optical mode distortion and higher order mode excitation when moving from the nominal waveguide width (i.e., 470 nm in this paper—see Section V) to a wider waveguide and vice versa [18]. In particular, the taper length should be long enough to avoid any optical transmission and mode distortion. Using Lumerical MODE [19], we simulated the fundamental mode transmission between two waveguides of different widths (e.g., R6 with six MZIs in Fig. 5). Accordingly, we consider the average observed non-uniform variations in a region to design region-based-tolerant MZIs by performing an exhaustive search for the MZI waveguide widths (using results in Fig. 3) while considering the area overhead in the MZIs imposed by adding the tapers. Here, we might have different waveguide widths (between the search range of 350 nm to 1200 nm—see Fig. 3) on each MZI arm.

As we will show in Section V, our region-based-tolerant and worst-case-tolerant MZI designs minimize optical phase noises in MZIs under different FPVs, hence they improve the network accuracy in SPNNs. Nevertheless, it is important to note that our device-level design optimization solutions proposed in this section do not aim at completely eliminating, if at all feasible, the impact of FPVs in SPNNs. That being said, our optimization will reduce the impact of FPVs in SPNNs sufficiently so that the overhead and complexity of dynamic calibration techniques (e.g., [6]) to eliminate the impact of such variations in SPNNs will be significantly reduced.

V. SIMULATION RESULTS AND DISCUSSIONS

FPVs lead to undesired optical phase noises, which, in turn, lead to faulty matrix-vector multiplication in the fully connected layers of an SPNN. The sensitivity of a standalone MZI to FPVs depends on its physical design parameters (see Fig. 3), among which only the MZI waveguide width can be freely altered during the design time. Leveraging realistic and correlated FPV maps developed based on [13] (see Fig. 2-top) and the proposed MZI design optimization in Section IV, we explore and optimize the nominal waveguide widths in the MZIs in SPNN case studies considered in this section, to improve their tolerance under different FPVs.

Prior to evaluating the impact of the proposed MZI optimization on SPNN accuracy under FPVs, let us explore whether such an optimization is independent of the model and the dataset. To examine this, Fig. 6(a) considers 100 randomly generated 16 × 16 unitary matrices—each of which belongs to a different weight matrix—and presents a box plot of the distribution of RV/D’s between the nominal and the deviated unitary matrices under FPVs, and for different region sizes. Recall from Section IV that to design the region-based-tolerant MZIs, we take the mean variations affecting a region on the FPV map with 1, 3, 6, or 12 MZIs into consideration (see Fig. 5). All the MZIs in a

![Fig. 4. Minimum taper length required to keep the optical transmission between two waveguides of different widths consistent (i.e., at 1 in the figure) and to avoid mode distortion. The inset zooms in the results for the taper length of 0–2 μm.](image)
particular region are replaced with the optimized region-based-tolerant MZI, designed using shallow-etched ridge waveguides (see Figs. 2 and 3(b)). As it is shown in Fig. 6(a), in all cases (i.e., R1–R12), the mean $RVD$ is significantly reduced when optimized MZIs are used. In particular, the interquartile ranges (IQRs) in the box plots are consistent among all the unitary matrices in a region: this shows that the proposed optimization is effective independent of the considered unitary matrix. Also, the mean $RVD$ decreases with decreasing the region size (i.e., when a designer has access to more detailed FPV data from the foundry), and it is the highest when MZIs are not optimized.

To explore the impact of our proposed MZI design optimization on SPNN accuracy, we consider a case study of two fully connected SPNNs with different footprint (see Table II) trained on the MNIST dataset. To compress the $28 \times 28 = 784$ dimensional feature vector in the MNIST dataset, we take the shifted fast Fourier transform of each image. The compressed 16-dimensional feature vector for Network-1 is then obtained by considering the values within the $4 \times 4$ region at the center of the frequency spectrum. Similarly, for the larger Network-2, we use a 64-dimensional feature vector by considering the $8 \times 8$ region at the center of the frequency spectrum. In Fig. 6(b), we show the linear correlation between the accuracy and the mean $RVD$, averaged over the 6 unitary matrices (two in each of the three OIUs) and normalized over the number of phase shifters in each network. Given such a strong linear correlation, our method should improve the accuracy of all SPNNs under FPVs, irrespective of the nominal phase angles.

By applying realistic and correlated FPV maps—generated using our prior work in [13] and considering parameters in Table I—to Network-1 and Network-2, Fig. 7 shows the inferencing accuracy in each network with conventional MZIs (No Optim.) and optimized region-based-tolerant MZIs, which can have different waveguide widths on each arm. For the conventional MZI, we designed an MZI using strip waveguides with $t = 220$ nm and $w = 470$ nm and variation-free $\approx 10$-nm-long DCs with a 200 nm gap, to obtain 50:50 splitting ratio in the MZI. In each plot in Fig. 7, we consider three different standard deviations for the waveguide width, SOI thickness, and slab thickness variations: $0.5 \times \sigma_w$, $1 \times \sigma_t$, and $2 \times \sigma_h$ in Table I. Moreover, the variation maps are generated for two correlation lengths of 100 μm and 1 mm. Considering Fig. 7(a)–(d), with no MZI optimization (i.e., No Optim.; red bars), the SPNN accuracy is always the least (e.g., 7.73% with $1 \times \sigma_t$ in Fig. 7(d)). Considering optimized region-based-tolerant MZI design with strip or shallow-etched ridge waveguides under all the variations and R1 (i.e., when

\begin{table}[h]
\centering
\begin{tabular}{|c|c|c|}
\hline
Model & Architecture & # PhS \\
\hline
Network-1 & FC(16,16)-SP-FC(16,16)-SP-FC(16,10)-LSM & 1380 \\
Network-2 & PC(64,64)-SP-PC(64,64)-SP-PC(64,10)-LSM & 20,580 \\
\hline
\end{tabular}
\end{table}
regions include a single MZI—see Fig. 5), the network accuracy in Fig. 7(a)–(d) is almost the same as the nominal accuracy. This is because each individual MZI has been optimized considering its exact FPV profile. This is for an ideal case when a designer has full access to variation data affecting each MZI in the network, so it may be impractical in most cases (R1 results are shown to indicate the efficiency of our optimization in such rare cases).

Considering more realistic and practical region sizes (R3, R6, and R12) and a region-based-tolerant MZI design using strip waveguides, Fig. 7(a)–(d) show the SPNN accuracy when considering (i) both waveguide width and SOI thickness variations (W+T; light green bars), and (ii) width variations only (W; magenta bars)—when SOI thickness variations are negligible, e.g., through SOI thickness uniformity improvement [20]. Observe that with (i), the optimized MZI can help retrieve some accuracy, which also decreases as the region size increases. But overall, the gain in accuracy in this case is small. This is due to the fact that optimized MZIs designed using strip waveguides are not sufficiently tolerant to thickness variations (see Fig. 3(a)). Considering (ii), the optimized MZI designed using strip waveguides achieves better accuracy improvements (and higher than (i)) in both the networks.

Fig. 7(a)–(d) also show the network accuracies with the optimized region-based-tolerant MZIs designed using shallow-etched ridge waveguides, under the presence of (i) waveguide width and SOI thickness variations (W+T; yellow bars), and (ii) all the variations (W+T+H; dark blue bars). Observe that with both (i) and (ii), the optimized MZIs using shallow-etched ridge waveguides perform much better compared to those using strip waveguides (i.e., ST (W+T)). Comparing (i) and (ii)—when considering slab thickness variations—the network accuracy is (slightly) lower in some cases. Nevertheless, considering the average difference between No Optim. accuracy and the accuracy obtained with (ii) for R1, R3, R6, and R12 (orange bars; second y-axis), using optimized MZIs designed with shallow-etched ridge waveguides under all the variations can significantly improve the network accuracy (e.g., by up to 72% in Network-2 with 0.5×σ). This shows that the proposed optimization can improve the resilience of SPNNs to FPVs.

Another observation is how the network accuracy in Fig. 7(a)–(d) seem to change across different region sizes and FPV correlation lengths. In fact, when the correlation length in variations is shorter, variations tend to “change more” within a region with a given size, and as the region size increases, they change even more across the region. This imposes a higher error for the region-based-tolerant MZIs designed per region. Therefore, the accuracy results are generally a bit lower in Fig. 7(a) and (c) compared to those in Fig. 7(b) and (d).

To assess SPNN accuracy using the optimized worst-case-tolerant MZI design (see Section IV), we consider, as an example, Network-2 with FPVs of different correlation lengths and standard deviations in Table I. In this experiment, we assume no a priori FPV knowledge and using strip waveguides in the design of the optimized worst-case-tolerant MZIs while considering waveguide width variations only. The worst-case-tolerant MZI is optimized by widening all the MZI arms together—i.e., MZI arms all have the same width after optimization—while considering the resulting area overhead due to the required waveguide tapers (see Figs. 2, 3(a), and 4). Results for this experiment (before and after the optimization) are shown in Table III for different area overhead. We observe that even with 1% area overhead, the accuracy improves. However, the improvements are significant only when the area overhead is greater than 8% for both the correlation lengths.

### VI. Conclusion

In this article, we have analyzed the impact of FPVs in the waveguide width, SOI thickness, and slab thickness on coherent SPNNs. In particular, we have modeled undesired optical phase noises due to such variations at the MZI device level, and how such phase noises contribute to the performance degradation, for which we have considered relative variation distance (RV D) at the network level in optical unitary multipliers built using MZIs. Furthermore, we have proposed physical-level design optimization solutions to enhance MZI device tolerance under correlated FPVs in SPNNs. Our simulation results for two SPNN case studies of different sizes and considering realistic FPV maps show that the proposed physical-level optimization can help significantly improve the SPNN inferencing accuracy. In addition, the results in this article indicate the importance of considering variations during the design-phase of SPNNs to facilitate the application of online and dynamic calibration mechanisms in these networks, which are often complex and power- and area-hungry.

### References

[1] F. P. Sunny, E. Taheri, M. Nikdast, and S. Pasricha, “A survey on silicon photonics for deep learning,” *ACM J. Emerg. Technol. Comput. Syst.*, vol. 17, no. 4, pp. 1–57, 2021.

[2] Q. Cheng, J. Kwon, M. Glick, M. Bahadori, L. P. Carloni, and K. Bergman, “Silicon photonics codesign for deep learning,” *Proc. IEEE*, vol. 108, no. 8, pp. 1261–1282, Aug. 2020.

[3] E. Cartlidge, “Optical neural networks,” *Opt Photon. News*, vol. 31, no. 6, pp. 32–39, 2020.

[4] S. Banerjee, M. Nikdast, and K. Chakraborty, “Characterizing coherent integrated photonic neural networks under imperfections,” *J. Lightw. Technol.*, early access, Aug. 2, 2022, doi: 10.1109/JLT.2022.3193658.

[5] Y. Zhu et al., “Countering variations and thermal effects for accurate optical neural networks,” in *Proc. IEEE/ACM Int. Conf. Comput. Aided Des.*, 2020, pp. 1–7.
[6] S. Bandyopadhyay et al., “Hardware error correction for programmable photonics,” *Optica*, vol. 8, no. 10, pp. 1247–1255, 2021.

[7] M. Y. et al., “Design of optical neural networks with component imprecisions,” *Opt. Exp.*, vol. 27, no. 10, pp. 14009–14029, 2019.

[8] M. Reck et al., “Experimental realization of any discrete unitary operator,” *Phys. Rev. Lett.*, vol. 73, pp. 58–61, 1994.

[9] W. R. Clements et al., “Optimal design for universal multiport interferometers,” *Optica*, vol. 3, no. 12, pp. 1460–1465, 2016.

[10] F. Shokraneh et al., “The diamond mesh, A phase-error- and loss-tolerant field-programmable MZI-based optical processor for optical neural networks,” *Opt. Exp.*, vol. 28, pp. 23495–23508, 2020.

[11] M. Nikdast, G. Nicolescu, J. Trajkovic, and O. Liboiron-Ladouceur, “Chip-scale silicon photonic interconnects: A formal study on fabrication non-uniformity,” *J. Lightw. Technol.*, vol. 34, no. 16, pp. 3682–3695, Aug. 2016.

[12] S. Banerjee et al., “Optimizing coherent integrated photonic neural networks under random uncertainties,” in *Proc. IEEE/OSA Opt. Fiber Comm. Conf. Exhib.*, 2021, pp. 1–3.

[13] A. Mirza, F. Sunny, P. Walsh, K. Hassan, S. Pasricha, and M. Nikdast, “Silicon photonic microring resonators: A comprehensive design-space exploration and optimization under fabrication-process variations,” in *IEEE Trans. Comput.-Aided Des. Integr. Circuits Syst.*, vol. 41, no. 10, pp. 3359–3372, Oct. 2022.

[14] F. Shokraneh, M. S. Nezami, and O. Liboiron-Ladouceur, “Theoretical and experimental analysis of a 4 × 4 reconfigurable MZI-based linear optical processor,” *J. Lightw. Technol.*, vol. 38, no. 6, pp. 1258–1267, Mar. 2020.

[15] M. M. P. Fard et al., “Experimental realization of arbitrary activation functions for optical neural networks,” *Opt. Exp.*, vol. 28, no. 8, pp. 12138–12148, 2020.

[16] Z. Lu et al., “Broadband silicon photonic directional coupler using asymmetric-waveguide based phase control,” *Opt. Exp.*, vol. 23, no. 3, pp. 3795–3808, 2015.

[17] G. Jiang, R. Chen, Q. Zhou, J. Yang, M. Wang, and X. Jiang, “Slab-modulated sidewall Bragg gratings in silicon-on-insulator ridge waveguides,” *IEEE Photon. Technol. Lett.*, vol. 23, no. 1, pp. 6–8, Jan. 2011.

[18] Y. Fu et al., “Efficient adiabatic silicon-on-insulator waveguide taper,” *Photon. Res.*, vol. 2, no. 3, pp. A41–A44, 2014.

[19] Lumerical, “Ansys lumerical.” Accessed: Apr. 2022. [Online]. Available: https://www.lumerical.com/products/

[20] S. K. Selvaraja et al., “SOI thickness uniformity improvement using corrective etching for silicon nano-photonic device,” in *Proc. IEEE Int. Conf. Group IV Photon.*, 2011, pp. 71–73.