Phase relations, and mechanical and electronic properties of nickel borides, carbides, and nitrides from \textit{ab initio} calculations†

Nursultan E. Sagatov, a*ab Aisulu U. Abuova, *b Dinara N. Sagatova, *c Pavel N. Gavryushkin, ac Fatima U. Abuova b and Konstantin D. Litasov d

Based on density functional theory and the crystal structure prediction methods, USPEX and AIRSS, stable intermediate compounds in the Ni–X (X = B, C, and N) systems and their structures were determined in the pressure range of 0–400 GPa. It was found that in the Ni–B system, in addition to the known ambient-pressure phases, the new nickel boride, Ni2B3–Immm, stabilizes above 202 GPa. In the Ni–C system, Ni3C–Pnma was shown to be the only stable nickel carbide which stabilizes above 53 GPa. In the Ni–N system, four new phases, Ni3N–R3, Ni5N–Cmcm, Ni7N3–Pmca, and NiN2–Pa3, were predicted. For the new predicted phases enriched by a light-element, Ni2B3–Immm and NiN2–Pa3, mechanical and electronic properties have been studied.

In the present paper, we investigate nickel borides, carbides, and nitrides under high pressure using modern \textit{ab initio} crystal structure prediction methods. The Ni–B system comprises several stable (NiB, Ni5B4, Ni3B, and NiB) and metastable (Ni23B6 and Ni3B4) crystalline compounds.15–19 Recently, Wang \textit{et al.}28 estimated the stability of ambient-pressure phases up to 100 GPa. They showed that Ni3B and Ni5B4 are unstable at pressures of 80–100 GPa, while Ni23B6 and NiB remains stable relative to the decomposition into isochemical mixture. In the Ni–C system, there is no stable intermediate compounds at ambient pressure. Recently, first stable nickel carbide, Ni3C with cementite-type structure, was synthesized in a diamond anvil cell (DAC) with laser heating through a direct reaction of a nickel powder with carbon at 184 GPa.25 In the Ni–N system, there is one known stable intermediate compound, Ni5N, at ambient pressure. Guillaume \textit{et al.}22 show that the ambient-pressure structure of Ni5N remains stable at 20 GPa. Recently, nickel pernitride (NiN2) with marcasite-type structure was synthesized at 40 GPa using a laser-heated DAC.23 It should be noted that no crystal structure prediction calculations were performed for these systems.

In our work, we search for new stable structures and study the phase stability, electronic and mechanical properties of the structures in the Ni–B, Ni–C, and Ni–N systems within first-principles calculations.

**Computational methods**

The search for new stable crystal structures of nickel borides, carbides, and nitrides were performed using evolutionary algorithms implemented in the USPEX package24–26 and the random sampling method implemented in the AIRSS software27–28 at 50, 100, 200, 300, and 400 GPa.

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* Sobolev Institute of Geology and Mineralogy, Siberian Branch of the Russian Academy of Sciences, Novosibirsk 630090, Russian Federation. E-mail: sagatino23@gmail.com
* L. N. Gumilyov Eurasian National University, Nur-Sultan 010008, Republic of Kazakhstan. E-mail: aisulu-us1980@yandex.ru
* Novosibirsk State University, Novosibirsk 630090, Russian Federation
*Vereshchagin Institute for High Pressure Physics, Russian Academy of Sciences, Troitsk, Moscow 108840, Russian Federation
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The search with USPEX was performed in variable composition mode with 6–36 atoms per unit cell. The size of the first generation in the calculations was equal to 65 structures. 60% of the structures with the lowest enthalpy were selected after the optimization and then used to produce the next generation. A new generation was produced as follows: 35% of all structures were generated by heredity, 20%—by atomic mutation, 10%—by lattice permutation, and 35%—randomly. In average, 36–40 generations have been produced and relaxed at each pressure. Using AIRSS, about 3500–4000 structures were randomly generated and optimized for compounds with fixed stoichiometries (NiXₙ; X = B, C, and N; n = 1–4) at each pressure, and those with the lowest enthalpy were selected.

The total energies and forces were calculated by solving the Schrödinger equation based on projector augmented plane-wave implementation of density functional theory (DFT) using the VASP package. Exchange correlation effects were treated in the generalized gradient approximation (GGA) with Perdew–Bürke–Ernzerhof scheme. In all crystal structure prediction calculations, medium-quality optimization was performed using the conjugate gradient. The medium quality settings were as follows: plane-wave cutoff energy—420 eV; Monkhorst-Pack k-point sampling grid of spacing—0.5 Å⁻¹; Gaussian smearing with parameter σ = 0.15 eV. The most promising predicted structures were then optimized at various pressures with higher accuracy: the cutoff energy—700 eV, k-point sampling grid of spacing—0.2 Å⁻¹, and σ = 0.05 eV.

To investigate mechanical properties, static elastic stiffness tensor (Cᵢⱼ) was calculated from the stress (σ) – strain (ε) relation σᵢⱼ = Cᵢⱼεᵢⱼ. To investigate electronic properties, in addition to the standard PBE, Dudarev’s GGA + U (ref. 32) and hybrid functional Heyd–Scuseria–Ernzerhof (HSE06) methods were used.

Structures were visualized in the VESTA package, and symmetries of the predicted structures were analyzed using the FindSym program.

### Results and discussion

#### Crystal structure prediction

The thermodynamic stability of the most favorable structures among the predicted ones was estimated using the convex hull diagrams, which give the convenient graphical representation of the thermodynamic stability of the predicted structures.

Crystal structure prediction calculations have been performed only for the intermediate compounds in the Ni–X (X = B, C, and N) systems. Crystal structure prediction calculations of pure Ni performed earlier by similar procedure showed its stability in Fm3m (fcc) structure up to 400 GPa. For pure boron, carbon, and nitrogen the following data have been used: B-R3m, -Pnma, -Cmca, and -P6₃/mcm; C-P6₃/mmc (graphite) and -Fd3m (diamond); N-P63, -P3, -Pba₂, and -I₄/mcm. The pressures of phase transition between listed polymorphic modifications are shown in Fig. 1b–3b.

![Fig. 1](image1.png)  
**Fig. 1** Thermodynamic convex hulls of Ni–B system at various pressures and zero temperature (a). Filled symbols denote stable structures, open symbols—metastable structures. Predicted pressure–composition phase diagram for the Ni–B system (b).

![Fig. 2](image2.png)  
**Fig. 2** Stability convex hulls of the Ni–C system at various pressures and zero temperature (a). Filled symbols denote stable structures, open symbols—metastable structures. Predicted pressure–composition phase diagram for the Ni–C system (b).
The Ni-B system. The Ni–B system is most rich in intermediate compounds with representatives in both Ni-rich and B-rich parts. In the Ni-rich part of the system four intermediate compounds, Ni3B, Ni2B, NiB2, and NiB, were revealed (Fig. 1a and S1†).

The Ni3B-Pnma phase was predicted at 50 GPa and it is stable from ambient pressure to 80 GPa. Above 80 GPa, Ni3B-Pnma decomposes into solid Ni and NiB2 (Fig. 1a and Table S1†). For Ni3B, the structure with the I4/mcm symmetry was predicted at 50, 100, 200, 300, and 400 GPa as the lowest-enthalpy one. This structure is stable in the entire considered pressure range without any phase transitions (Fig. 1b). The Ni2B-Pnma was predicted at 50 GPa as the most energetically favorable. It is stable from 0 to 86 GPa, and above this pressure decomposes into the mixture of (NiB2 + NiB). The NiB was predicted in the Cmcm structure, which stabilizes at a pressure of 2 GPa and remains stable up to 400 GPa.

In the B-rich part of the Ni–B system, we predicted a new compound that was not previously known, namely Ni3B3 (Fig. 1 and Table S1†). This compound was found at 200, 300, and 400 GPa with Immm symmetry. Above 202 GPa, the enthalpy of Ni3B3-Immm becomes lower than that of mechanical mixture of (2NiB + B), indicating that the Ni3B3-Immm structure stabilizes relative to the decomposition reaction Ni3B3 ↵ 2NiB + B (Fig. S2a†).

All these Ni–B structures are dynamically stable, which were confirmed by the calculated phonon spectra (Fig. S3†).

Although the search of the structures stable at ambient pressure has not been performed, the comparison of the found structures with the known ones show, that all the structures stable at ambient pressure, which were observed in experiments, were revealed in our calculations at 50 GPa. In addition, we have considered Ni12B6, which was observed in experiments, and Ni3B3 (ref. 19) suggested theoretically. According to our calculations, Ni12B6 is metastable at 0 GPa. This is in agreement with experimental results. For Ni4B3, we considered two structures: Ni2B3-P63mc suggested in ref. 19 and Ni2B3-Pbca constructed from Fe2C3-Pbca, replacing Fe by Ni and C by B atoms. These structures are very similar and their distinguishing by X-ray powder diffraction technique is problematic. Our calculation showed that Ni4B3-Pbca has lower enthalpy than Ni2B3-P63mc in the pressure range of 0–155 GPa (Fig. S4a†). Based on this, we suggest that most likely the Pbca phase was formed in experiments of Hofmann et al., ¹⁹ than hexagonal phase suggested in that work. With pressure increase, Ni12B6 and Ni3B3 becomes even less favorable (Fig. 1).

The Ni–C system. Crystal structure prediction calculations have revealed one stable compound in the Ni-C system, Ni3C (Fig. 2 and Table S2†). It stabilizes in cementite-type orthorhombic structure Pnma, which is in agreement with recent experimental results, ²¹ where Ni3C-Pnma was synthesized at 184 GPa using DAC technique. According to the obtained results, Ni3C-Pnma becomes energetically more favorable than (3Ni + C) above 53 GPa (Fig. 2b). There are no imaginary modes in the phonon spectra of Ni3C-Pnma (Fig. S5†), indicating that it is dynamically stable.

In addition, we constructed Ni-analogs of stable iron carbide structures Fe7C3-P63mc, ²⁴ Fe2C3-Pbca, ²⁴ and Fe2C3-Pnma, ²⁴ by replacing all iron atoms by nickel ones. According to the obtained results, constructed Ni3C-Pbca is more favorable than constructed Ni2C3-P63mc (Fig. S4b†). Nevertheless, Ni3C-Pbca is unstable relative to the mechanical mixtures of (Ni + C) or (NiC + C) in the entire considered pressure range (Fig. 2). Similar calculations for Ni3C-Pnma showed that this compound is not stable, and the energy difference reaches 0.4 eV f.u.⁻¹.

Crystal structure predictions in C-rich part of the Ni–C system were performed for fixed stoichiometries, namely NiC4, NiC3, NiC2, and NiC. For all these compounds the lowest-enthalpy structures are unstable and decompose into isochemical mixture in the entire considered pressure range.

The Ni–N system. The preliminary results of crystal structure prediction calculations in the Ni–N system in the pressure range of 0–300 GPa were reported as the conference proceedings. ⁴⁶ Here, we extended the studied pressure range up to 400 GPa, found new structures based on the analogy with carbides and reanalyze obtained phase relations. ⁴⁶ In addition to crystal structure predictions, we constructed two structural models, Ni2N2-Pbca and Ni3N2-P63mc, from known iron carbides Fe2C3-Pbca and Fe2N2-P63mc by replacing all Fe by Ni atoms. The volume, lattice parameters, and atomic coordinates were optimized at various pressures for the resulting models.

Performed crystal structure prediction calculations show that there are two stable intermediate compounds in the Ni–N system at 50 GPa, Ni1N and NiN2 (Fig. 3 and Table S3†). Both...
predicted Ni$_3$N-P6$_2$22 and NiN$_2$-Pnnm were observed earlier in high-pressure experiments. According to obtained results, Ni$_3$N-P6$_2$22 is stable in the pressure range of 4–96 GPa. Below 4 GPa, Ni$_3$N-P6$_2$22 is unstable relative to mechanical mixture of (3Ni + N), and above 96 GPa it is unstable relative to mechanical mixture of [Ni$_3$B$_3$ + Ni]. NiN$_2$-Pnnm stabilizes above 14 GPa. At 100 GPa, there are three stable nickel nitrides, predicted Ni$_6$N-P3 and NiN$_2$-Pa3, and constructed Ni$_3$N-Pbca. Ni$_6$N-P3 stabilizes at 98 GPa, and above 114 GPa decomposes into solid Ni and Ni$_2$N$_2$. Constructed Ni$_3$N-Pbca is more favorable than constructed Ni$_3$N-P6$_2$mc in the entire pressure range (Fig. S2†). Ni$_3$N-Pbca become stable relative to Ni$_3$N + Ni$_2$N$_2$ above 93 GPa and stable at least up to 400 GPa. Ni$_2$N$_2$ undergoes a phase transition from known NiN$_2$-Pnnm to the new predicted NiN$_2$-Pa3 at 96 GPa. NiN$_2$-Pa3 is stable at least up to 400 GPa. At 400 GPa, in addition to Ni$_3$N-Pbca and NiN$_2$-Pa3, one new nitride Ni$_3$N-Cmcm was predicted. It stabilizes relative to the mechanical mixture of [Ni$_3$N$_3$ + Ni] above 331 GPa. Phonon calculations show that the predicted nickel nitrides are dynamically stable; i.e., they have no imaginary modes in the phonon spectra (Fig. S6†).

Crystal chemistry of the predicted structures

Structural data of the new predicted compounds are given in Table 1.

The most crystallographically interesting structures are the structures of compounds enriched by the light element. We have revealed two such compounds, NiN$_2$ and Ni$_2$B$_3$. Two found structures of Ni$_3$N, NiN$_2$-Pa3, corresponds to the well-known structural type of pyrite with N–N molecular units.

The crystal structure of Ni$_2$B$_3$ is also characterised by the covalently bonded units. In this case, these are infinite (001) layers of boron atoms. In the layer each B atom is connected to four other B atoms forming the net with hexagonal and rhombic loops (Fig. 4a). The B–B bond distance within the layer is 1.67 Å at 250 GPa. The elemental boron is characterised by nearly the same B–B distance, 1.7 Å. The covalent nature of B–B bond in Ni$_2$B$_3$ structure is also confirmed by the electronic density distribution (Fig. S8†). Boron layers alternates with layers of Ni atoms, forming rectangular centered network (Fig. 4b). Each Ni is coordinated by 10 boron atoms, the six of which are from the top, and four are from the bottom layer. The performed analysis of the structures with the similar stoichiometry has revealed another structure, W$_2$CoB$_2$, belonging to the described structural type.47,48

Graphical representation of the Ni$_3$N-P3, Ni$_3$N-Cmcm, Ni$_2$N$_2$-Pbca, and NiN$_2$-Pa3 crystal structures are shown in Fig. S7†.

Elastic properties

Transition metal borides, carbides, and nitrides often exhibit extraordinary mechanical properties. Elastic constants are helpful to understand the mechanical properties and provide

Table 1  Structural data of structures introduced in the present work

| Compound | $P$ (GPa) | Space group | Lattice parameters ($\text{Å}$, deg) | Atom | Coordinates |
|----------|----------|-------------|-----------------------------------|------|-------------|
| Ni$_2$B$_3$ | 250 | Immm (171) | $a = 2.695$  $b = 3.357$  $c = 5.622$  | Ni1 | 0.000  0.000  0.804 |
|           |         |             | $\alpha = 90.00$  $\beta = 90.00$  $\gamma = 90.00$ | B1 | 0.500  0.000  0.500 |
| Ni$_2$N | 100 | R3 (148) | $a = 4.058$  $b = 4.058$  $c = 11.653$ | Ni1 | 0.316  0.331  0.246 |
|           |         |             | $\alpha = 90.00$  $\beta = 90.00$  $\gamma = 120.00$ | N1 | 0.000  0.000  0.000 |
| Ni$_2$N$_2$ | 100 | Pbca (161) | $a = 4.131$  $b = 12.671$  $c = 10.978$ | Ni1 | -0.080  0.534  0.686 |
|           |         |             | $\alpha = 90.00$  $\beta = 90.00$  $\gamma = 90.00$ | Ni2 | 0.744  0.878  0.585 |

Fig. 4 Crystal structure of Ni$_2$B$_3$-Immm in two projections.
useful information to estimate the hardness of a structure. To that end, we investigate the elastic properties of new predicted compounds enriched by a light-element, NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm, within the strain–stress method in combination with the DFT. Calculated elastic constants of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are listed in Table 2. From Table 2 and Fig. 5, it could be seen that the elastic constants $C_{ij}$ for these structures are all positive, increase monotonically with increasing pressure, and the whole set of the elastic constants of the structures satisfy the mechanical stability criteria. This means that NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are mechanically stable in the considered pressure range.

Calculated bulk modulus $B$ of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are ~87% and ~82% of that estimated for diamond, respectively, while shear modulus $G$ of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are ~35% and ~38% of that estimated for diamond, respectively (Table 3). Predicted NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are ductile materials since their $B/G$ values (Table 3) are greater than 1.75.

In order to investigate the hardness of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm at various pressures we employ empirical Chen$^{35}$ and Tian et al.$^{32}$ models. The Vickers hardness of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are estimated to be 19 GPa and 12 GPa, respectively. The hardness of NiN$_2$-Pa3 decreases with increasing pressure, while that of Ni$_2$B$_3$-Immm increases with increasing pressure. The estimated $H_v$ of both predicted structures are below minimal criteria of 20 GPa of hardness. Therefore, NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm are not classified as a hard materials.

**Electronic structure**

To investigate electronic characteristics of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm, we calculate their total/partial electronic density of states (DOS) (Fig. 6) and band structures (Fig. 7).

The Fermi level of the considered phases is set to zero. The result shows that Ni$_2$B$_3$-Immm exhibit metallic features. From calculated partial DOS (Fig. 6a), we noticed that Ni d orbital has hybridization with B p orbital. The lower bands in the valence band region are mainly contributed by B 2p states, while near the top of valence bands main contributors are Fe 3d states. The Fermi level are mostly occupied by the Fe 3d electrons, with some contribution of B 2p electrons. On the other hand, NiN$_2$-Pa3 is found to be a semiconductor with the indirect band gap of 1.297 eV. The electronic states at the lower energy range of

### Table 2 Calculated elastic constants $C_{ij}$ (GPa) of NiN$_2$-Pa3 and Ni$_2$B$_3$-Immm at different pressures

| Phase      | $P$ (GPa) | $C_{11}$ | $C_{12}$ | $C_{13}$ | $C_{22}$ | $C_{23}$ | $C_{33}$ | $C_{44}$ | $C_{55}$ | $C_{66}$ |
|------------|-----------|----------|----------|----------|----------|----------|----------|----------|----------|----------|
| NiN$_2$-Pa3| 100       | 1311.6   | 326.4    | 326.2    | 1311.5   | 326.5    | 1311.6   | 204.8    | 204.8    | 204.8    |
|            | 200       | 1887.7   | 609.7    | 609.5    | 1887.7   | 609.7    | 1887.7   | 243.5    | 243.5    | 243.5    |
|            | 300       | 2363.3   | 891.5    | 891.3    | 2363.3   | 891.5    | 2363.3   | 263.1    | 263.1    | 263.1    |
|            | 400       | 2787.9   | 1175.5   | 1175.2   | 2787.9   | 1175.5   | 2787.9   | 272.8    | 272.8    | 272.8    |
| Ni$_2$B$_3$-Immm | 200       | 1372.3   | 831.2    | 724.6    | 1354.5   | 774.6    | 1429.9   | 305.3    | 169.8    | 493.4    |
|            | 300       | 1773.5   | 1114.3   | 973.2    | 1750.0   | 1893.5   | 404.0    | 239.6    | 643.6    |          |
|            | 400       | 2153.0   | 1389.9   | 1218.3   | 2126.0   | 1307.1   | 2336.1   | 486.2    | 297.1    | 780.4    |

Fig. 5 Calculated elastic constants of Ni$_2$B$_3$-Immm (a) and NiN$_2$-Pa3 (b) as functions of pressure.

| Phase      | $P$ (GPa) | $B$    | $G$    | $B/G$   | $H_v^{\text{Chen}}$ | $H_v^{\text{Thm}}$ |
|------------|-----------|-------|-------|--------|---------------------|---------------------|
| NiN$_2$-Pa3| 100       | 654.8 | 293.6 | 2.23   | 18.7               | 20.6               |
|            | 200       | 1035.7| 362.7 | 2.86   | 15.4              | 18.1               |
|            | 300       | 1382.1| 403.2 | 3.43   | 12.8              | 15.8               |
|            | 400       | 1712.9| 428.6 | 4.00   | 10.7              | 13.9               |
| Ni$_2$B$_3$-Immm | 200       | 979.7 | 297.5 | 3.29   | 10.9              | 13.4               |
|            | 300       | 1298.1| 388.1 | 3.34   | 12.9              | 15.9               |
|            | 400       | 1604.8| 467.7 | 3.43   | 14.2              | 17.6               |

Diamond$^{33}$

| Phase      | $P$ (GPa) | $B$    | $G$    | $B/G$   | $H_v^{\text{Chen}}$ | $H_v^{\text{Thm}}$ |
|------------|-----------|-------|-------|--------|---------------------|---------------------|
|            | 100       | 838   | 848   |        | —                   | —                   |
|            | 200       | 1188  | 1023  |        | —                   | —                   |
|            | 300       | 1521  | 1168  |        | —                   | —                   |
|            | 400       | 1839  | 1295  |        | —                   | —                   |

Fig. 6 Total and partial DOS calculated for Ni$_2$B$_3$-Immm (a) and NiN$_2$-Pa3 (b).
the valence bands are mainly contributed by N 2p orbitals, while Ni 3d orbitals dominate at the top of valence bands. The main contributors in the conduction band region are unoccupied Fe 3d states, with some contribution of N 2p states.

It is well-known that GGA often underestimates the width of band gap. For this reason, and also to avoid misleading, we performed GGA + \( U \) (with \( U = 2, 4 \) eV) and hybrid (HSE06) calculations. The obtained results showed that PBE, GGA + \( U \), and HSE06 give qualitatively the same DOS and band structures (Fig. S9 and S10†). This fact indicates that the results obtained by standard PBE calculations are valid. For NiN\(_2\)-Pa\(^3\), the band gap obtained by GGA + \( U \) with \( U = 2 \) eV was equal to 1.524 eV, by GGA + \( U \) with \( U = 4 - 1.001 \) eV, by HSE06 – 2.237 eV (Fig. S10b†).

To elucidate chemical bond properties of Ni\(_2\)B\(_3\)-Immm and NiN\(_2\)-Pa\(^3\), electron localization function (ELF) diagrams were constructed (Fig. 8). The ELFs of the structures are derived from various planes and different colours indicate diverse degrees of electron localization: the red means that the electron is totally localized at the realm and blue means the opposite situation, namely, the electron is not localized at all. According to the obtained result, in Ni\(_2\)B\(_3\)-Immm, we could observe that the boron atoms form a mesh structure. In NiN\(_2\)-Pa\(^3\), the nitrogen atoms form N\(_2\) dimers. In both structures, valence electrons are not localized around the Ni atoms. These data suggest that the chemical bonds among nickel mainly exhibit ionic characteristics and boron/nitrogen atoms form covalent chemical bonds, which seemingly do not strongly interact with nickel atoms.

**Conflicts of interest**

There are no conflicts to declare.

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