A novel memcapacitor and its application in a chaotic circuit

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Received: 8 February 2021 / Accepted: 10 June 2021 / Published online: 28 June 2021
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Abstract In this paper, a novel memcapacitor is designed by the SBT (Sr$_{0.95}$Ba$_{0.05}$TiO$_3$) memristor and two capacitors. A fifth-order memcapacitor and memristor chaotic circuit is proposed. The stability of the equilibrium point of the system is analyzed theoretically. Lyapunov exponents spectra, bifurcation diagrams, poincaré maps and phase diagrams are used to analyze the dynamic behaviors of the system. The results show that under different initial values and parameters, the system produces rich dynamic behaviors such as stable points, limit cycles, chaos and so on. Specially, coexisting attractors, transient chaos, and steady-state chaos accompanied by burst period phenomena are also produced in the system. The proposed memcapacitor-based circuit expands the research methods of memcapacitor for application in chaotic circuits.

Keywords Memcapacitor · Memristor · Chaotic circuit · Multistability

1 Introduction

Chua predicted the existence of memristor based on mathematical relations between pairs of fundamental circuit variables in 1971 [1]. HP (Hewlett–Packard) laboratory made the first solid-state thin-film memristor in 2008 [2]. A memristor is a nonlinear two-terminal circuit element reflecting the relationship between magnetic flux and charge ($\varphi - q$). Chua et al. further proposed memcapacitor based on memristor, which further expanded the concept of memory elements [3,4]. Memcapacitor is a passive two-terminal electronic device described by the nonlinear constitutive relationship of magnetic flux and the time-domain integral of electric charge ($\varphi - \sigma$). Memcapacitor or memristor can easily generate high-frequency chaotic oscillation signals in the circuit because of the characteristic of nonlinearity, and many chaotic circuits introduced them to obtain chaos [5–9]. The chaotic circuits can generate some unique phenomena, including transient chaos [10–13], steady-state chaos accompanied by burst period phenomena [14], and multistable phenomena [15–19].

The behaviors of different memcapacitor models [20–26] were explored through previous researches on memcapacitor. Reference [20] described a synthesis of mutators which can transform the emulated memristor into memcapacitor and meminductor. Reference [21] converted a digital memristor into memcapacitor by virtue of the voltage following characteristics of operational amplifiers, which complicated the memcapacitor model. Reference [22] presented a memcapacitor simulator based on an light-dependent resistor (LDR) memristor and constructed a memcapacitor model. Refer-
ence [23] designed a floating memcapacitor emulator, the circuit structure was simple and can be widely used in circuit design. Recently, researches based on memcapacitor have become a focus [27–29]. In Reference [27], a fractional-order memcapacitor model was proposed and a chaotic oscillator based on the model was investigated. Reference [28] proposed a logarithmic charge-controlled memcapacitor model and verified it non-volatility and switching features. These simulating memcapacitors were mostly based on complex conversion circuits, which were prone to errors, and it is not easy to the analysis of the characteristics of the memcapacitor. Therefore, it is necessary to find a new way to study the application of memcapacitor.

The memcapacitor device can be implemented by appending a memristor with a MIM (metal–insulator–metal) capacitor in Reference [30]. As the resistance of memristor changes, the capacitance of memcapacitor changes under external excitation. This structure has a potential for the artificial neural networks and chaotic circuits. The memcapacitor device also has nonlinear characteristics and low power consumption, and its structure is simpler than some simulating memcapacitors, so it has potential application in the integrability and scalability of CMOS technology. Based on the method of constructing physical memcapacitor proposed by Reference [30], this paper uses the physical SBT memristor [31, 32] to design a memcapacitor, which can be directly used as a circuit element to generates chaos in the circuit. This paper constructs a fifth-order chaotic circuit composed of a memcapacitor and a memristor.

The paper is organized as follows. In Sect. 2, the dynamical modeling of the fifth-order memcapacitor and memristor chaotic circuit is introduced and its corresponding plane equilibrium and stability are analyzed. In Sect. 3, the influences of initial states and circuit parameters on system dynamic behaviors are studied. The conclusions are drawn in Sect. 4.

2 Dynamical modeling of the chaotic circuit based on memcapacitor and SBT memristor

2.1 The chaotic circuit based on memcapacitor and SBT memristor

In the previous researches, the mathematical model of the memristor had been proposed [31, 32]. According to the experimental measurement of the memristor, the flux-controlled model of the memristor was obtained as follows:

\[
\begin{align*}
    & i(t) = (A + B |\varphi(t)|) u(t) \\
    & \frac{d\varphi(t)}{dt} = u(t)
\end{align*}
\]

where \( A = 0.0676 \), and \( B = 0.3682 \).

The fifth-order chaotic circuit based on the novel memcapacitor and the SBT memristor is shown in Fig. 1. The fifth-order chaotic circuit consists of a novel memcapacitor, a SBT memristor \( W_2 \), a resistor \( R \), a negative conductance-\( G \), and an inductance \( L \). According to Reference [30], the novel memcapacitor is composed of a memristor \( W_1 \) and two capacitors \((C_1 \) and \( C_2)\), as shown in the dotted line. Its structure is simpler than previous memcapacitors [20–22], and it can be directly applied to the design of chaotic circuits without establishing a model of memcapacitor.

According to Kirchhoff’s circuit laws, the current \( i_L \) of inductor \( L \), the voltage \( u_1 \) of capacitor \( C_1 \), the voltage \( u_2 \) of capacitor \( C_2 \), the magnetic flux \( \varphi_1 \) of memristor \( W_1 \), and the magnetic flux \( \varphi_2 \) of memristor \( W_2 \) are selected as state variables, the state equations of the system are as follows:

\[
\begin{align*}
    & \frac{di_L(t)}{dt} = \frac{1}{L} \left[ u_1(t) + u_2(t) - Ri_L(t) - \frac{i_L(t)}{W_2} \right] \\
    & \frac{du_1(t)}{dt} = \frac{1}{C_1} \left[ G(u_1(t) + u_2(t)) - i_L(t) \right] \\
    & \frac{du_2(t)}{dt} = \frac{1}{C_2} \left[ G(u_1(t) + u_2(t)) - i_L(t) - W_1 u_2(t) \right] \\
    & \frac{d\varphi_1(t)}{dt} = u_2(t) \\
    & \frac{d\varphi_2(t)}{dt} = \frac{i_L(t)}{W_2}
\end{align*}
\]

Let \( x = i_L(t), y = u_1(t), z = u_2(t), w = \varphi_1(t), v = \varphi_2(t), a = 1/C_1, b = 1/C_2, c = 1/L, r = \ldots \)
Table 1  The system parameters for the chaotic attractor

| Parameters | $a$ | $b$ | $c$ | $r$ | $g$ |
|------------|-----|-----|-----|-----|-----|
| Values     | 3.39| 1.95| 9.36| 0.42| 1.42|

Fig. 2  The typical chaotic attractor of fifth-order chaotic circuit

When $R$, and $g = G$, the dimensionless equations of the system are as follows:

$$
\begin{align*}
\dot{x} &= c[y + z - rx - \frac{x}{W_2}] \\
\dot{y} &= a[g(y + z) - x] \\
\dot{z} &= b[g(y + z) - x - W_1z] \\
\dot{w} &= z \\
\dot{v} &= \frac{z}{W_2}
\end{align*}
$$

where $W_1 = A + B|w|$, $W_2 = A + B|v|$.

2.2 Typical chaotic attractors

When the parameters are set as shown in Table 1, the complex dynamic behaviors occur in the system. The system generates a double-scroll attractor (see Fig. 2). By Jacobi matrix method, the five Lyapunov exponents are calculated as $\text{LE}_1 = 0.1446$, $\text{LE}_2 = 0.01265$, $\text{LE}_3 = 0.0091$, $\text{LE}_4 = -0.2864$, and $\text{LE}_5 = -6.917$. The sum of the Lyapunov exponents is negative, which means that the system is chaotic.

2.3 Plane equilibrium and stability distribution

Let $\dot{x} = \dot{y} = \dot{z} = \dot{w} = \dot{v} = 0$ in Eq. (3), and the circuit parameters are set as shown in Table 1, the equilibrium point of the system can be obtained as:

$$
A = \{(x, y, z, w, v) \mid x = y = z = 0, w = m, v = n\}
$$

where $m$ and $n$ are real constants. That is, the points on the $w-v$ plane are the equilibrium point of the system (3). The Jacobian matrix $J$ of Eq. (3) can be expressed as:

$$
J = \begin{bmatrix}
-cr - \frac{c}{W_2} & c & c & 0 & 0 \\
-a & ag & ag & 0 & 0 \\
-b & bg & bg & -bW_1 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 \\
\end{bmatrix}
$$

The characteristic equation of the equilibrium point $A$ can be given as:

$$
\det(\lambda E - J) = \lambda^2(\lambda^3 + a_1\lambda^2 + a_2\lambda + a_3) = 0
$$

where $a_1, a_2, a_3$ are as follows:

$$
\begin{align*}
a_1 &= bW_1 - bg - ag + cr + \frac{c}{W_2} \\
a_2 &= (bcr - abg)W_1 + \frac{c}{W_2} - acrg - bcrg + ac + bc \\
a_3 &= (abc - abcr)W_1 - \frac{abcgW_1}{W_2}
\end{align*}
$$

According to the Routh–Hurwitz stability criterion, if all the nonzero eigenvalues of Eq. (6) are negative, the system is stable:

$$
a_1 > 0, a_3 > 0, a_1a_2 - a_3 > 0
$$

If $n = 0$, $m$ is a variable parameter, the conditions $a_1 > 0$, and $a_3 > 0$ in Eq. (8) cannot be satisfied simultaneously. If $m = 0$, $n$ is a variable parameter, Eq. (8) has no solution. Obviously, the equilibrium point $A$ is always unstable. No matter where the system (3) starts, the system tends to limit cycles, chaos, or infinite divergence.

3 The complex dynamic behaviors of the system

3.1 Chaos oscillations dependent on the initial states

3.1.1 Influences of initial conditions $x(0), y(0), z(0), w(0)$, and $v(0)$ on system chaos oscillations

When the dynamic behaviors of the system change with the initial states, the system always oscillates in chaos.
The circuit parameters are taken as shown in Table 1, and the initial values of the system are assigned as $(0.001, 0, 0, 0, 0)$. The Lyapunov exponents spectra and bifurcation diagrams varying with the initial states $x(0)$, $y(0)$, $z(0)$, $w(0)$, and $v(0)$ are shown in Figs. 3, 4, 5, 6 and 7. In order to better observe the Lyapunov exponents of the chaos, the fifth Lyapunov exponential curves are not drawn in Figs. 3, 4, 5, 6 and 7. The system remains chaos oscillations with the variation of initial states.

3.1.2 Multistability depending on the initial condition $z(0)$

Multistability is a common characteristic in chaotic systems. Figure 8 shows the coexisting chaotic attractors and the corresponding Poincaré maps. Figure 8a, b shows the coexisting single-scroll chaotic attractors in detail and the corresponding Poincaré map, where the blue attractor starts from the initial conditions of $(0.001, 0, 7, 0, 0)$ and the red one starts from $(0.001, 0, -7, 0, 0)$. Figure 8c, d shows the coexisting double-scroll chaotic attractors in detail and the corresponding Poincaré map, where the blue attractor starts from the initial conditions of $(0.001, 0, 1.5, 0, 0)$ and the red one starts from $(0.001, 0, -1.5, 0, 0)$. Figure 8e, f shows the coexisting single-scroll chaotic attractor and double-scroll chaotic attractor in detail and the corresponding Poincaré map, where the blue attractor starts from the initial conditions of $(0.001, 0, 3, 0, 0)$ and the red one starts from $(0.001, 0, -3, 0, 0)$.
3.1.3 Multistability depending on the initial conditions $x(0)$, $y(0)$, $w(0)$, and $v(0)$

Other initial conditions can also produce coexisting chaotic attractors. The circuit parameters are taken as shown in Table 1, and the phase diagrams and the Poincaré maps of the multistability dependent on the initial conditions $x(0)$, $y(0)$, $w(0)$, and $v(0)$ are shown in Figs. 9, 10, 11 and 12, the results are summarized in Table 2.
Fig. 11 The multistability of the circuit varying with initial conditions $w(0)$. a the phase diagram when initial conditions are (0.001, 0, 0, 3, 0) and (0.001, 0, 0, −5, 0); b the Poincaré map when initial conditions are (0.001, 0, 0, 3, 0) and (0.001, 0, 0, −5, 0); c the phase diagram when initial conditions are (0.001, 0, 0, 1.5, 0) and (0.001, 0, 0, −2, 0); d the Poincaré map initial conditions are (0.001, 0, 0, 1.5, 0) and (0.001, 0, 0, −2, 0).

Fig. 12 The multistability of the circuit varying with initial conditions $v(0)$. a the phase diagram when $v(0) = \pm 5$; b the Poincaré map when $v(0) = \pm 5$.

3.2 Dynamic behaviors dependent on system parameters

3.2.1 Influences of parameter $a$ on system dynamic behaviors

When parameter $a$ is in the range of [1.70, 3.80] and initial values are (0.001, 0, 0, 0, 0), the Lyapunov exponents spectrum and bifurcation diagram can be obtained, as shown in Fig. 13. In order to better observe the Lyapunov exponents of the chaos, the fifth Lyapunov exponential curve is not drawn. When $a$ is in the range of [1.70, 1.81], the system is stable; when $a$ is in the range of [1.82, 3.67], the system exhibits chaotic attractors; when $a$ is in the range of [3.68, 3.80], the system exhibits limit cycles. The dynamical evolution process of the system with the change of parameter $a$ is shown in Fig. 14, where the values of $a$ are 3.22, 3.39, and 3.70, respectively.

When parameter $a$ is in the range of [1.70, 3.80], and others are set as shown in Table 1. Dynamical behaviors with coexisting bifurcation diagrams are pre-

| Figure | Initial conditions | Coexisting attractors |
|--------|-------------------|----------------------|
| Figure 9a, b | (±5, 0, 0, 0, 0) | Single-scroll |
| Figure 9c, d | (±1.8, 0, 0, 0, 0) | Double-scroll |
| Figure 10a, b | (0.001, ±4, 0, 0, 0) | Single-scroll |
| Figure 10c, d | (0.001, ±0.4, 0, 0, 0) | Double-scroll |
| Figure 11a, b | (0.001, 0, 0, 3, 0) and (0.001, 0, 0, −5, 0) | Single-scroll |
| Figure 11c, d | (0.001, 0, 0, 1.5, 0) and (0.001, 0, 0, −2, 0) | Double-scroll |
| Figure 12a, b | (0.001, 0, 0, 0, ±5) | Single-scroll |
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Fig. 15 The coexistence bifurcation diagrams of parameter $a$ with the initial values $(0.001, 0, 0, 0, 0)$ and $(0, 0.001, 0, 0, 0)$

Fig. 16 The coexisting signal-scroll chaotic attractors plotted for $a = 3.22$. (a) the phase diagram when initial values are $(0.001, 0, 0, 0, 0)$ and $(0, 0.001, 0, 0, 0)$; (b) the Poincaré map when initial values are $(0.001, 0, 0, 0, 0)$ and $(0, 0.001, 0, 0, 0)$

Fig. 17 a Lyapunov exponents spectrum and b bifurcation diagram varying with the system parameter $c$

Fig. 18 The coexistence bifurcation diagrams of parameter $c$ with the initial values $(0.001, 0, 0, 0, 0)$ and $(-0.001, 0, 0, 0, 0)$

3.2.2 Influences of parameter $c$ on system dynamic behaviors

Parameter $c$ is set in the range of $[9.00, 14.00]$, and other parameters are set as shown in Table 1. When the initial values are $(0.001, 0, 0, 0, 0)$, the Lyapunov exponents spectrum and bifurcation diagram can be obtained, as shown in Fig. 17. In order to better observe the Lyapunov exponents of the chaos, the fifth Lyapunov exponential curve is not drawn. When parameter $c$ is in the range of $[9.00, 13.37]$, the system is chaotic, while in the range of $[13.38, 14.00]$, the system is stable.

Dynamical behaviors with coexisting bifurcation diagrams are presented in Fig. 18, where the trajectories colored in blue start from the initial conditions $(0.001, 0, 0, 0, 0)$ and the corresponding Poincaré map are shown in Fig. 19, where the blue attractor starts from the initial conditions $(0.001, 0, 0, 0, 0)$ and the red one starts from $(-0.001, 0, 0, 0, 0)$. 
Fig. 19 The coexisting signal-scroll chaotic attractors plotted for \( c = 9.76 \) and initial conditions are \((\pm 0.001, 0, 0, 0, 0)\). \( a \) the phase diagram; \( b \) the Poincaré map

Fig. 20 \( a \) Lyapunov exponents spectrum and \( b \) bifurcation diagram varying with the system parameter \( b \)

Fig. 21 \( a \) Lyapunov exponents spectrum and \( b \) bifurcation diagram varying with the system parameter \( r \)

Fig. 22 \( a \) Lyapunov exponents spectrum and \( b \) bifurcation diagram varying with the system parameter \( g \)

Table 3 The dynamics with the variation in system parameters \( b, r, \) and \( g \)

| System parameters | Interval     | Dynamics                      |
|-------------------|--------------|-------------------------------|
| \( b \)           | [0.00, 0.32] | Stable point                  |
|                   | [0.33, 0.64] | Single-scroll attractor       |
|                   | [0.65, 2.00] | Double-scroll attractor       |
| \( r \)           | [0.30, 0.35] | Limit cycle                   |
|                   | [0.36, 0.42] | Double-scroll attractor       |
|                   | [0.43, 0.50] | Single-scroll attractor       |
| \( g \)           | [1.00, 1.14] | Stable point                  |
|                   | [1.15, 1.40] | Single-scroll attractor       |
|                   | [1.41, 1.50] | Double-scroll attractor       |

3.3 Transient chaos

The initial values of the system are set as \((0.001, 0, 0, 0, 0)\), other parameters remain unchanged. When \( b = 1.97 \), the chaotic phenomenon of the system is transient chaos in the finite time range. The phenomenon of transient chaos accompanied by boundary crisis is often encountered in dynamic systems. As shown in Fig. 23, with the state variable \( v \) changing in the time range \([20 \text{ s}, 360 \text{ s}]\) and \([380 \text{ s}, 1000 \text{ s}]\), the system produces two different dynamic phenomena. When \( t \in [20 \text{ s}, 360 \text{ s}] \), the system behaves as single-scroll attractors. As time goes by, when \( t \in [380 \text{ s}, 1000 \text{ s}] \), the system behaves as double-scroll attractors.

3.4 Steady-state chaos accompanied by burst period phenomenon

When \( b = 0.65 \), the system has a strange phenomenon of steady-state chaos accompanied by burst period. Figure 24 shows the steady-state chaotic attractors with
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4 Conclusion

In this paper, a novel memcapacitor based on the SBT memristor is proposed. It has a simple structure and can be directly used in the design of chaotic circuits. Then, a fifth-order chaotic circuit based on this memcapacitor is designed. The characteristics of the circuit are analyzed by using the dynamics methods. With the changing of initial conditions \((x(0), y(0), z(0), w(0), v(0))\), the system is always chaotic and produces a wealth of coexisting attractors. With the changing of system parameters, the system generates complex dynamic behaviors such as transient chaos, steady-state chaos accompanied by burst period. The proposed memcapacitor and memristor chaotic circuit in this paper enriches the application of memcapacitor and memristor in high-order circuits and expands the research methods of memcapacitor in chaotic circuits.

Acknowledgements

This work was supported by the National Natural Science Foundation of China (Grant Nos. 61703246, 61703247), the Qingdao Science and Technology Plan Project (Grant Nos. 19-6-2-2-cg), and the Elite Project of Shandong University of Science and Technology.

Declaration

Conflicts of interest

The authors declare that they have no conflict of interest.

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Fig. 23 The phenomenon of transient chaos when \(b = 1.97\). a time-domain trajectory in the interval of \([0 \, s, 1000 \, s]\); b the phase diagram when \(t \in [20 \, s, 360 \, s]\) and \(t \in [380 \, s, 1000 \, s]\)

Fig. 24 The phenomenon of steady-state chaos accompanied by burst period when \(b = 0.65\)

Fig. 25 The coexisting phenomenon of steady-state chaos accompanied by burst period plotted for \(b = 0.65\) and initial conditions are \((\pm 0.001, 0, 0, 0, 0)\). a the phase diagram; b the Poincaré map

period 8 in different planes, that is, a period 8 orbits coexist in chaos. This phenomenon is sensitive to the initial values of the system. When \(b = 0.65\), the coexisting phenomenon of steady-state chaos accompanied by burst period in detail and the corresponding Poincaré map are shown in Fig. 25, where the blue attractor starts from the initial conditions of \((0.001, 0, 0, 0, 0)\) and the red one starts from \((-0.001, 0, 0, 0, 0)\).
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