Defects-assisted piezoelectric response in liquid exfoliated MoS$_2$ nanosheets

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Abstract
MoS$_2$ is an intrinsic piezoelectric material which offers applications such as energy harvesting, sensors, actuators, flexible electronics, energy storage and more. Surprisingly, there are not any suitable, yet economical methods that can produce quality nanosheets of MoS$_2$ in large quantities, hence limiting the possibility of commercialisation of its applications. Here, we demonstrate controlled synthesis of highly crystalline MoS$_2$ nanosheets via liquid phase exfoliation of bulk MoS$_2$, following which we report piezoelectric response from the exfoliated nanosheets. The method of piezo force microscopy was employed to explore the piezo response in mono, bi, tri and multilayers of MoS$_2$ nanosheets. The effective piezoelectric coefficient of MoS$_2$ varies from 9.6 to 25.14 pm V$^{-1}$. We attribute piezoelectric response in MoS$_2$ nanosheets to the defects formed in it during the synthesis procedure. The presence of defects is confirmed by x-ray photoelectron spectroscopy.

Keywords: 2D materials, MoS$_2$, piezoelectricity

(Some figures may appear in colour only in the online journal)

1. Introduction
Piezoelectricity is the charge that originated when certain materials are placed under stress. The stretching or compressing of the substance induces electricity and allows the facile turning of mechanical energy into electrical and vice versa which offers an immense application for sensors, energy harvesting, flexible electronics, energy storage, and actuators. A great deal of research is being carried out in generating clean energy due to obvious reasons [1, 2]. It is very essential to find sustainable and environment friendly non-conventional forms of energy. To substitute conventional fuels, one way is to harvest electricity from human surroundings. The piezoelectric systems could prove extremely useful in such a scenario. Piezoelectricity at micro and nanoscale can be effectively engineered by necessary surface modifications including the introduction of adatoms and defects in such materials leading to surface piezoelectricity. Vast amount of research has been reported on piezoelectricity in various inorganic and organic based materials including composites [3–5]. However, the above-mentioned piezoelectric materials have some drawbacks due to their lower piezoelectric properties, non-biodegradability, rigidity, toxic nature, non-ecofriendliness, and non-biocompatibility. In search of efficient piezoelectric material, 2D materials are emerging as a potential alternative as theoretically predicted and some of them have been confirmed experimentally [6–8]. The transition metal dichalcogenides (TMDCs) in particular have been explored increasingly due to their non-centrosymmetric crystal structure and large surface area to thickness ratio [9–11]. As a typical TMDC material, MoS$_2$ is of great interest due to its novel properties such as tunable band gap, large in-plane mobility, and mechanical stability [12–15]. Each layer of MoS$_2$ is made up of Mo atoms sandwiched between two layers of hexagonally close-packed S atoms, and the weak van der Waals forces tie the sandwiched adjacent layers [16]. Ultrathin MoS$_2$ layers can be prepared through various techniques such as hydrothermal method, liquid phase exfoliation (LPE), chemical vapor deposition, mechanical exfoliation, and so on [17–20]. Among these techniques, LPE method looks promising for scalable production of MoS$_2$ because of its low cost,
relatively high output and simplicity. But controlled synthesis of crystalline MoS₂ with desired structure and morphology is still a challenge.

In case of MoS₂ it has been observed that the odd layered MoS₂ is highly piezoelectric due to non-centrosymmetric nature, with monolayer showing the strong piezoelectricity [9, 21]. The piezoelectric effect disappears in bulk MoS₂ due to the opposite orientations of the adjacent layers. Most of these demonstrations of piezoelectric effect have been made on mechanically exfoliated odd number of nanosheets which pose limitations for large scale applications. As an alternative, the technique of LPE along with fabrication techniques like solution casting, spray coating, or inkjet printing can be used for scalable fabrication of electronic devices. In this direction, some theoretical research on piezoresponse in multilayer 2D materials has been reported but intensive experimental studies are still lacking [9, 11, 22]. To the best of our knowledge, no experimental report is available on the piezoelectric response from the liquid phase exfoliated few layers of MoS₂. In this work, we propose that 2D MoS₂ layers synthesized via modified LPE method contain defects that can act as prime spots for trapping charges and display piezo signals.

2. Experimental details

The MoS₂ nanosheets were prepared by modified LPE method. In detail, 2g of MoS₂ powder (<6 μm, 98%, Sigma) was dispersed in 60 ml water and probe sonicated for 1 h to get a homogeneous mixture. The obtained mixture was kept in a hydrothermal autoclave reactor and subjected to continuous stirring for 3 h at 150 °C. The pre-treated MoS₂ was diluted in IPA/water and further exfoliated using a probe sonicator for about 5 h. The obtained dispersion was centrifuged at 5000 rpm for 30 min to collect the supernatant. The collected supernatant was further diluted before coating on a substrate to avoid nanosheet agglomeration. Few trials of the spin coating were made to get un-agglomerated coating of nanosheets on ITO substrate for all PFM measurements.

3. Characterization

The surface morphology of MoS₂ nanosheets was observed and analyzed by scanning electron microscope coupled with energy dispersive x-ray (SEM-EDX) spectroscopy (Ultra55 FE-SEM Karl Zeiss). All the SEM images were taken under the operating voltage of 5 kV in InLens mode. The qualitative analysis was performed using energy dispersing x-ray spectroscopy (EDS) at electron high tension voltage of 15 kV. The Raman spectra were taken at room temperature using Horiba LabRAM HR Raman spectrometer. Instrument resolution was ~0.5 cm⁻¹ by using 1800 groove mm⁻¹ grating. A linearly polarized laser light (532 nm) was used for the measurements keeping the laser power below 1 mW. Structural analysis and mapping were carried out using a transmission electron microscope (TEM TITEN Themis) conducted at accelerating voltage of 200 kV. For TEM analysis, diluted MoS₂ dispersion was drop cast over a carbon-coated copper grid, followed by drying under IR bulb. The PFM measurements were carried out using an atomic force microscope (AFM), Park system (Park AFM NX 20). A conducting cantilever (Multi75E-G) having a force constant of 3 N m⁻¹ and resonant frequency 75 kHz was employed in all our PFM measurements.

4. Results and discussion

The morphology of MoS₂ layers was characterized by TEM as shown in figure 1(a), which suggests the presence of very thin sheets. A clear view of lattice fringes having the lattice spacing of the order of 0.27 nm which corresponds to the (100) plane of MoS₂ nanosheets, can be observed in figure 1(b) [10]. The inset of figure 1(b) shows the SEAD pattern of MoS₂ confirming the crystallinity of the exfoliated nanosheets. Raman spectra in figure 1(c) show the peak difference of ~25 cm⁻¹ between the characteristic peaks of MoS₂ confirming the presence of a few layers [23]. Also, the photoluminescence spectra of MoS₂ layers shown in figure 1(d) reveal that the band gap value of synthesized MoS₂ nanosheets is around 1.75 eV.

Figures 2(a) and (b) shows SEM image of exfoliated nanosheets and the corresponding lateral area distribution of the flakes respectively. The average flake area is found to be 0.4663 μm². Figure 2(c) shows the EDS spectra of MoS₂ nanosheets and the spectra confirm that the as-synthesized MoS₂ nanostructures contain Mo and S elements along with oxygen. The AFM image of MoS₂ nanosheets is shown in figure 2(d). The average roughness (Ra) and root mean square roughness (Rq) of MoS₂ layers is 1.2 nm and 3.1 nm respectively. The thickness distribution of exfoliated nanosheets, analyzed using AFM is shown in figure 2 (e). The average flake thickness is around 2.96 nm. This suggested that bulk MoS₂ particles are successfully exfoliated and significantly fragmented.

To quantify the piezoelectric response at the nanoscale, the PFM based investigations were carried out using AFM in contact mode. For local piezoelectric study, common AFM provides an ideal platform, due to high vertical resolution and high localization of electric field at the junction between the metalized tip and the surface. The PFM uses the basic principle of inverse piezoelectric effect wherein AC field is applied onto the sample surface through AFM probe keeping the bottom electrode grounded, which results in the surface deformation of the sample. The amplitude/phase of this precise deformation could be measured to characterize the strength/direction of the dipole. The tip-sample voltage is adjusted with a periodic tip bias

\[ V_{tip} = V_{DC} + V_{AC} \cos(\omega t) \]  

(1)

here the drive frequency \( \omega \) is chosen to be well above the feedback bandwidth. This drive is responsible for generating an oscillating electric field below the tip that can induce localized deformation in the MoS₂ surface.
During the measurement, the cantilever oscillation frequency was directed close to its resonant frequency with $V_{\text{bias}} = 6$ V. The drive frequency of 50 kHz was chosen to optimize the signal stability and signal-to-noise ratio, as there was less than 20% variation in PFM amplitude signal in the frequency range of 20–60 kHz.

The PFM measurements were carried out on a few layers of MoS$_2$, where the nanosheets were not agglomerated. Throughout the PFM measurements, the cantilever tip (Bruker SCM-PIT) of AFM acts as a top electrode and ITO coated substrate as the bottom electrode. Figures 3(a)–(c) shows the topography, amplitude, and phase image of typical MoS$_2$ layers deposited on to ITO substrate. The deflection sensitivity of the system (48 V $\mu$m$^{-1}$) was determined using force-distance (F-D) curve (figure 3(d)) by indenting the cantilever tip on a hard surface of sapphire. To get electrostatic free piezoresponse, the voltage was applied during PFM measurements and a stiffer (2.945 N m$^{-1}$) cantilever was used throughout the measurement. The comparison of the input voltage with periodic surface vibrations using a lock-in amplifier was used to quantify the PFM response in terms of its amplitude ($A$) and phase ($\varphi$) signals. The characteristic butterfly loop in the PFM amplitude and the hysteresis loop in the phase signal as a function of applied voltage for a few layers of MoS$_2$ are also shown in figures 3(e) and (f) respectively. To determine the value of the piezoelectric constant, the piezoresponse was measured by different driving voltages repeatedly. The change in PFM amplitude indicates mechanical responses with respect to the electrical signal of $V_{\text{tip-cons}}$ confirming the piezoelectric effect from MoS$_2$ nanosheets. To quantify the effect, the linear piezoelectricity can be determined by the following equation $^{[24]}$

$$A = d_{33} V_{\text{AC}},$$

where $A$ is the PFM amplitude, $V_{\text{AC}}$ is the drive ac bias voltage and $d_{33}$ is the effective piezoelectric coefficient. The measurements were conducted in the strong-indentation regime, where the piezo-response is dominated by the $d_{33}$ rather than $d_{31}$ of the material. Here $d_{33}$ is the induced strain in direction 3 per unit electric field applied in direction 3 whereas $d_{31}$ is the induced strain in direction 1 per unit electric field applied in direction 3. From the plot of piezo-response amplitude versus $V_{\text{bias}}$, taken on the MoS$_2$ layers...
we extracted the effective piezoelectric coefficient value ($d_{33}$) to be 9.6 pm V$^{-1}$. Due to the small-scale effect, the classical piezoelectricity caused by the non-centrosymmetric crystal structure will not be dominant in our case [25]. To confirm that this response was not coming only because of the classical piezoelectricity of MoS$_2$ in odd number of layers, the piezo response of MoS$_2$ having different thicknesses (monolayer, bilayer, and trilayer) was investigated by PFM under similar imaging parameters. All the PFM measurements were carried out in the
ramping voltage range of $-10\,\text{V}$ to $+10\,\text{V}$ for MoS$_2$ having a different thicknesses as shown in figure 4. The AFM topography of the MoS$_2$ sample, consisting of mono, bi, and tri-layer of MoS$_2$ is marked as 1, 2, and 3 respectively in figure 4(a). The PFM amplitude signal as a function of ramped $V_{\text{bias}}$ at the frequency of 50 kHz for respective mono-, bi- and tri-layer of MoS$_2$ are shown in figures 4(b)–(d). Different AC voltages were applied from the AFM tip to calculate the piezoelectric coefficient, which reveals a linear relationship between the applied electric field and mechanical deformation. As expected, a strong piezoelectric response was observed from monolayer MoS$_2$. From the piezo-response amplitude versus $V_{\text{bias}}$ curves taken on the spots marked by the plus signs (figure 4(a)), we extracted the effective piezoelectric coefficient ($d_{33}$) value of 10.91 pm V$^{-1}$, 15.75 pm V$^{-1}$, and 25.14 pm V$^{-1}$ for monolayer, bilayer, and tri-layer respectively. The decrease in $d_{33}$ value in the case of a few layers may be due to the screening effect induced by free carriers. The material with a higher piezoelectric coefficient produces more electricity for a given strain thus improving the efficiency of the piezoelectric device.

In general, the piezoelectric effect appears when external mechanical stress breaks the atomic symmetry of a material, which causes the accumulation of electric charges. In the case of even layered MoS$_2$, the piezoelectricity from alternating layers is canceled out. Therefore, the piezo response has been observed in odd number of layers of MoS$_2$ [9]. The piezoelectric response from MoS$_2$ in this work is probably coming because of defects as observed through EDS analysis. There is a possibility of trapping and stabilization of charges around the defects in the material as reported by A. Apte et al [26] that can induce a piezoelectric response. The synthesis procedure in this work gives rise to oxidation which suggests an efficient route to artificially engineer piezoelectricity in TMDC.

To further verify the presence of defects, XPS measurements have been performed. In the full survey spectra shown in figure 5(a), the characteristic peaks of Mo, S, O, and Si (due to substrate) atoms were observed. In the image, except for C1s peak, no other impurity peak can be found. The corresponding binding energy plots for Mo 3d, S 2p, and O 1s are presented in figures 5(b)–(d) respectively. All spectra were referenced by adventitious carbon 1s peak at 284.8 eV (not shown). In the Mo 3d spectra of as-prepared MoS$_2$ layers, the characteristic peaks at 229.25 eV and 231.41 eV

![Figure 4](image-url)
belong to the Mo$^{4+}$ 3d$_{5/2}$ and Mo$^{4+}$ 3d$_{3/2}$ components of MoS$_2$, respectively [27, 28]. The peak at 227.27 eV arises due to S 2s component of MoS$_2$ [27]. The binding energies of Mo 3d and S 2p regions reveal that the few layer MoS$_2$ has a strong trigonal prismatic structure. The peak at binding energies 232.59 eV and 235.6 eV are ascribed to Mo$^{6+}$ revealing that the Mo edges in MoS$_2$ layers are slightly oxidized during the transition from the Mo$^{4+}$ state to the Mo$^{6+}$ state which corresponds to the Mo(VI) species (Mo$^{6+}$ 3d$_{5/2}$ and Mo$^{6+}$ 3d$_{3/2}$) in MoO$_3$ [29, 30]. Therefore, peaks attribute to the formation of Mo–O suggesting some degree of MoO$_x$-like structures. Here in MoS$_2$, oxygen can be attached as substitutional atoms at sulfur sites or as atoms bound to Mo atoms at plane edges. Additionally, there are two strong peaks at 162.07 eV and 163.3 eV which are assigned to S 2p$_{3/2}$ and S 2p$_{1/2}$ binding energies for S$^{2-}$. The peak at 169.2 eV could be ascribed to the existence of an S–O bond which indicates partial oxidation of the S edges in MoS$_2$ layers. The oxidation percentage of Mo and S edges was 12.04% and 27.13% respectively. The charges and dipoles are trapped at the defect sites of MoS$_2$ flakes. These induced defects act as nucleation and pinning centers for moving domain walls. It has been observed earlier that the motion of the domain wall is a significant contributor to the piezoelectric coefficient [31]. Hence, we believe, the defects induced in MoS$_2$ flakes during the synthesis process result into the piezoelectric response in the few layers of MoS$_2$.

5. Summary

In summary, a strong piezoelectric response in modified liquid exfoliated MoS$_2$ nanosheets has been observed. The strong piezoelectric coefficient in MoS$_2$ is attributed to the formation of defects during its synthesis process which is confirmed by XPS analysis. The defects promote the creation of a local asymmetric structure which gives rise to strong piezo-response. This work suggests an efficient route to artificially engineer piezoelectricity in MoS$_2$. Given the material’s properties and piezoelectric response, the liquid exfoliated MoS$_2$ could be considered a promising candidate for piezotronics devices and energy harvesting technology.

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Data availability statement

The data that support the findings of this study are available upon reasonable request from the authors.

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