Study of $e^+e^- \rightarrow D^+D^-\pi^+\pi^-$ at center-of-mass energies from 4.36 to 4.60 GeV

BESIII Collaboration

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We report a study of the $e^+e^- \rightarrow D^+D^-\pi^+\pi^-$ process using $e^+e^-$ collision data samples with an integrated luminosity of 2.5 fb$^{-1}$ at center-of-mass energies from 4.36 to 4.60 GeV, collected with the BESIII detector at the BEPCII storage ring. The $D_1(2420)^+$ is observed in the $D^+\pi^-\pi^-$ mass spectrum. The mass and width of the $D_1(2420)^+$ are measured to be $(2427.2 \pm 1.0_{\text{stat}} \pm 1.2_{\text{sys}})$ MeV/c$^2$ and $(23.2 \pm 2.3_{\text{stat}} \pm 2.3_{\text{sys}})$ MeV, respectively. In addition, the Born cross sections of the $e^+e^- \rightarrow D_1(2420)^+D^- + c.c. \rightarrow D^+D^-\pi^+\pi^-$ and $e^+e^- \rightarrow \psi(3770)\pi^+\pi^- \rightarrow D^+D^-\pi^+\pi^-$ processes are measured as a function of the center-of-mass energy.

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Table 1

The numbers relevant to the Born cross section measurements, where the first uncertainties are statistical, the second are independent systematic uncertainties, and the third are common systematics. The index of 1 represents the process $e^+e^+ \rightarrow D_1(2420)^+D^- + c.c. \rightarrow D^+\pi^-\pi^-\pi^+$ while the index of 2 represents the process $e^+e^+ \rightarrow \psi(3770)\pi^+\pi^- \rightarrow D^+D^-\pi^+\pi^-$. The upper limits correspond to the 90% confidence level. The symbol $\sigma$ refers to the statistical significance.

| $E_{cm}$ (MeV) | $\mathcal{L}$ (pb$^{-1}$) | $\sigma_{int}^B$ | $\epsilon_1$ (%) | $1 + \delta_{s1}^{\text{stat}}$ | $\sigma_1$ (pb) | $S_1$ |
|----------------|--------------------------|----------------|-----------------|----------------|--------------|------|
| 4358.3         | 543.9                    | 0.010 ± 0.010  | 23.90           | 0.795           | 1.051        | 39.8 5.3 ± 6.2 ± 3.5 | 7.9σ |
| 4378.4         | 55.6                     | 125 ± 28       | 23.20           | 0.821           | 1.051        | 59.8 3.3 ± 3.9 ± 5.3 | 4.9σ |
| 4415.6         | 1090.7                   | 2454 ± 111     | 22.56           | 0.820           | 1.053        | 61.9 2.8 ± 3.9 ± 5.5 | 24.9σ |
| 4467.1         | 111.1                    | 100 ± 28       | 20.92           | 0.904           | 1.055        | 24.1 6.6 ± 5.6 ± 2.1 | 39σ  |
| 4527.1         | 112.1                    | 122 ± 24       | 19.27           | 0.935           | 1.055        | 30.5 3.0 ± 3.0 ± 2.7 | 5.8σ |
| 4574.5         | 48.9                     | 24 ± 15 (± 43) | 18.22           | 1.029           | 1.055        | 13.2 8.3 ± 2.1 ± 1.2 (<23.7) | 1.7σ |
| 4599.5         | 586.9                    | 572 ± 56       | 17.92           | 1.075           | 1.055        | 25.6 2.5 ± 1.2 ± 2.3 | 11.7σ |

1. Introduction

Recent discoveries of charmonium-like states that do not fit naturally with the predicted charmonium states in the quark model have stirred up great experimental and theoretical interests [1–5]. Among these so-called XYZ states, the observations of the Y(4260) [6] and Zc(4430) [7] states have drawn special attention, and stimulated extensive discussions on their structures. Some calculations indicate that the Y(4260) is possibly a $D_1(2420)^+\bar{D}$ molecular state, while the Zc(4430) is possibly a $D_1(2420)^+\bar{D}$ molecular state [8–11]. Hence, more studies on the properties of the involved $D_1(2420)^+$, such as mass and width, are helpful to better understand the nature of these exotic candidate states.

The lightest charmonium state above the $\bar{D}D$ threshold is the $\psi(3770)$ resonance, which is considered to have the quantum numbers of $1^3D_1$ [12,13]. Its spin-triplet partner $1^3D_2$ candidate, $X(3823)$, has been observed in the process $e^+e^- \rightarrow X(3823)\pi^+\pi^-\pi^+$ at BESIII [14]. Analogously, it is interesting to study the production of the $\psi(3770)$ in the process $e^+e^- \rightarrow \psi(3770)\pi^+\pi^-\pi^+$ [15], which is observed at $\sqrt{s}=4415.6$ MeV at BESIII [16]. More precise measurements at different energy points are desired, as it provides an important way to investigate the intrinsic nature of the Y(4360) and $\psi(4415)$ by studying the transitions between these charmonium-like states, such as Y(4360) $\rightarrow \psi(3770)\pi^+\pi^-\pi^+$ and $\psi(4415) \rightarrow \psi(3770)\pi^+\pi^+$.

In this analysis, we study the process $e^+e^- \rightarrow D^+\bar{D}^0\pi^+\pi^-$ at the center-of-mass (c.m.) energies, from 4358.3 to 4599.5 MeV, as listed in Table 1. Compared to the process $e^+e^- \rightarrow D^0\bar{D}^0\pi^+\pi^-$, this final state has the advantage of being free from $D^*$ intermediate states, which greatly simplifies the analysis. We reconstruct the $D^*$ via its high branching fraction decay $K^+\pi^-\pi^+$ and adopt a recoil-mass technique to identify the $D^*$ and related resonant states. Unless explicitly mentioned otherwise, inclusion of charge conjugate mode is implied throughout the context. Clear signals of the $D_1(2420)^+$ and $\psi(3770)$ are extracted in this data set via their decays to $D^+\pi^+\pi^-\pi^+$ and $D^+\pi^-\pi^-$, respectively. The resonance parameters of the $D_1(2420)^+$ are measured. Additionally, the Born cross sections of $e^+e^- \rightarrow D_1(2420)^+D^- + c.c. \rightarrow D^+\pi^-\pi^-\pi^+$ and $e^+e^- \rightarrow \psi(3770)\pi^+\pi^- \rightarrow D^+D^-\pi^+\pi^-$ are measured at each $E_{cm}$.

2. The experiment and data sets

The BESIII detector is a magneto-spectrometer [17] located at the Beijing Electron Positron Collider (BEPCII) [18]. The cylindrical core of the BESIII detector consists of a helium-based multilayer drift chamber (MDC), a plastic scintillator time-of-flight system (TOF), and a CsI(Tl) electromagnetic calorimeter (EMC), which are all enclosed in a superconducting solenoidal magnet providing a 1.0T magnetic field. The solenoid is supported by an octagonal flux-return yoke with resistive plate counter muon identifier modules interleaved with steel. The acceptance of charged particles and photons is 93% over 4$\pi$ solid angle. The charged-particle momentum resolution at 1 GeV/c is 0.5%, and the $dE/dx$ resolution is 6% for the electrons from Bhabha scattering. The EMC measures photon energies with a resolution of 2.5% (5%) at 1 GeV in the barrel (end cap) region. The time resolution of the TOF barrel part is 68 ps, while that of the end cap part is 110 ps.

The $E_{cm}$ of the seven data sets are measured using di-muon events [19], and the corresponding luminosities are measured with large-angle Bhabha scattering events [21]. To optimize selection criteria, estimate the detection efficiency, and understand background contributions, we simulate the $e^+e^-$ annihilation processes with the KKMC [22] generator, which takes into account continuum processes, initial state radiation (ISR), and inclusive $D^{(*)}_{s1}$ production. The known decay rates are taken from the Particle Data Group (PDG) [13], and the decays are modeled with EVTGEN [23]. The remaining decays are simulated with the LUNDCHARM package [24]. The four-body process $e^+e^- \rightarrow D^+\bar{D}^0\pi^+\pi^-$ is generated considering the intermediate resonances $e^+e^- \rightarrow D_1(2420)^+D^-\pi^+$ assuming the relative orbital angular momentum of $D_1(2420)^+D^-$ in s-wave, and $e^+e^- \rightarrow \psi(3770)\pi^+\pi^-\pi^+$ assuming $\psi(3770)\pi^+\pi^-\pi^+$ uniformly distributed in momentum phase space, along with the subsequent decays $D_1(2420)^+ \rightarrow D^+\pi^+\pi^-$ and $\psi(3770) \rightarrow D^+\pi^-\pi^-$, respectively. We simulate one million events for each process at different $E_{cm}$. All simulated Monte Carlo (MC) events are processed in a geant4-based [25] software package, taking into account detector geometry and response.
3. Event selection and data analysis

3.1. Event selections

To reconstruct the $D^+$ meson, charged track candidates for one $K^-$ and two $\pi^+$ in the MDC are selected. For each track, the polar angle $\theta$ defined with respect to the $e^+e^-$ beam is required to satisfy $|\cos \theta| < 0.93$. The closest approach to the $e^+e^-$ interaction point is required to be within $\pm 10$ cm along the beam direction and within $\pm 1$ cm in the plane perpendicular to the beam direction. A track is identified as a $\pi(K)$ when the PID probabilities satisfy $P(\pi) > P(K) > P(\pi)$, according to the information of $dE/dx$ and TOF. We reconstruct $D^+$ candidates by considering all possible combinations of the charged tracks, which are required to originate from a common vertex. The quality of the vertex fit is required to satisfy $\chi^2_{\text{VF}} < 100$. We constrain the reconstructed $D^+$ mass with a kinematic fit to the nominal $D^+$ mass [13], and require the fit quality $\chi^2_{\text{VF}} < 20$. We then require the presence of one additional $\pi^+\pi^-$ pair, with neither track used in the reconstructed $D^+$. The identification of the signal process $e^+e^- \rightarrow D^+D^-\pi^+\pi^-$ is based on the recoil mass spectra of $D^+\pi^-\pi^-$, $RM(D^+\pi^+\pi^-)$, which are shown in Fig. 1. The rate of multiple candidates per event is about 10%, and is corrected for via the MC efficiency.

The peaks observed at 1.87 GeV/c² correspond to the $D^-$ meson signals. They are consistent with the MC simulations of the $D^+D^-\pi^+\pi^-$ final state. The background contributions are due to random combinations of charged tracks. We further restrict the candidate events to the region $1.855 < RM(D^+\pi^+\pi^-) < 1.882$ GeV/c², and plot the recoil mass of the $D^+$, $RM(D^+)$, as shown in Fig. 2. Enhancements around the $D_1(2420)\pi$ nominal mass are clearly visible. We take the events with $RM(D^+\pi^-\pi^-)$ in the sideband regions of (1.786, 1.840) GeV/c² and (1.897, 1.951) GeV/c² which are illustrated in Fig. 1, as samples representing the combinatorial background contributions in the distributions of $RM(D^+)$. This approach has been verified using the corresponding distributions of the background contributions from the inclusive MC samples. It is found that the sideband samples correctly reproduce the background in the signal region of $RM(D^+\pi^+\pi^-)$. Besides the contributions from $D_1(2420)\pi^+$, there is a clear excess of the data over background contributions from the sideband at high $RM(D^+)$ mass. It is consistent with from the process $e^+e^- \rightarrow \psi(3770)\pi^+\pi^- \rightarrow D^+D^-\pi^+\pi^-$. 

3.2. Signal extraction

The 2-dimensional distributions of $M(D^+\pi^+\pi^-)$ versus $RM(D^+)$ for the $D_1(2420)\pi$ are shown in Fig. 3. The vertical band corresponds to the $D_1(2420)\pi^-\pi^+$ signal and the horizontal band corresponds to the $D_1(2420)^+\pi^-\pi^+$. The projection to the $RM(D^+)$ axis (Fig. 2) consists of a prominent $D_1(2420)^+$ peak and a corresponding broad bump. The contributions of $D_1(2420)^+\pi^+$ and $\psi(3770)\pi^+\pi^-$ in the selected data are determined using fits to the $RM(D^+)$ one-dimensional distribution. The shape of this distribution is described using templates obtained from the signal MC simulation. In order to perform a likelihood scan of the resonance parameters, we generate a series of $D_1(2420)^+$ signal MC with different values of mass and width, and smear these template shapes with a Gaussian function to take into account the resolution difference between data and MC simulations. The width of the Gaussian function is fixed to the difference of resolution in $RM(D^+)$ for the control sample of $e^+e^- \rightarrow D^+D^-\pi^+\pi^-$. The signal shape for the mode $\psi(3770)\pi^+\pi^-\pi^-$ is obtained from the MC simulation, where the resonance parameters of the $\psi(3770)$ are taken from the PDG [13]. The relativistic Breit-Wigner function [13] is used to model the resonance lineshape of the $\psi(3770)$ and $D_1(2420)^+\pi^-\pi^+$.

A simultaneous unbinned maximum likelihood fit to the data samples is performed at three high luminosity energy points of $\sqrt{s} = 4383.3, 4415.6$ and 4599.5 MeV, with the resonance parameters of the $D_1(2420)^+$ in common for all fits. The shapes and magnitudes of the combinatorial backgrounds are fixed according to the sample of the sideband events in $RM(D^+\pi^-\pi^-)$, while the magnitudes of the $D_1(2420)^+\pi^+$ and $\psi(3770)\pi^+\pi^-\pi^-$ are the free parameters of the fit. The sum of the fitting components is shown in Fig. 2. We obtain the mass and width of the $D_1(2420)^+$ to be $2427.2 \pm 1.0$ MeV/c² and $(23.2 \pm 2.3)$ MeV, respectively. The signal yields are also measured, as listed in Ta-
Table 1. Here, the contribution of the non-resonant four-body process $e^+e^- \rightarrow D^+D^-\pi^+\pi^-$ is neglected in the fit, as an alternative fit including this process gives its size consistent with zero.

In addition, we analyze the data samples at $E_{c.m.} = 4487.4, 4467.1, 4527.1$ and $4574.5$ MeV with relatively low luminosities. We apply the same strategy to extract the signal yields of the $D_1(2420)^+D^-$ and $\psi(3770)\pi^+\pi^-$, except that we fix the resonance parameters for the $D_1(2420)^+$ according to the aforementioned fit results.

### 3.3 Cross section measurement

The Born cross section is calculated with

$$\sigma_i = \frac{\eta_{i,\text{sig}}}{2LB\varepsilon_1(1 + \delta_i^{\text{rad}})^2} \frac{1}{|1 - |t|^2]}, \quad \text{(1)}$$

where index $i$ denotes the respective signal process, $\eta_{i,\text{sig}}$ is the observed signal yield, $L$ is the integrated luminosity, $B$ is the branching fraction $B(D^+ \rightarrow K^-\pi^+\pi^-) = (9.38 \pm 0.16)\%$ [$15$], $\varepsilon_1$ is the detection efficiency, $(1 + \delta_i^{\text{rad}})$ is the radiative correction factor which is obtained from a QED calculation using the line shape of the data cross section of signal process as input in an iterative procedure, and $\frac{1}{|1 - |t|^2}$ is the vacuum polarization factor [26]. The trigger efficiencies for the two processes are $100\%$, as there are at least 5 charged tracks detected [27]. The processes $e^+e^- \rightarrow D_1(2420)^+D^- + c.c. \rightarrow D^+D^-\pi^+\pi^-$ and $e^+e^- \rightarrow \psi(3770)\pi^+\pi^- \rightarrow D^+D^-\pi^+\pi^-$ are denoted with index $i = 1$ and $i = 2$, respectively. The calculated Born cross sections are given in Table 1 and plotted in Fig. 4. We evaluate the statistical significance by the ratio of the maximum likelihood value and the likelihood value for a fit with a null-signal hypothesis. For the energy points with low statistical significances, we determine the upper limits for the cross sections which are calculated by using the signal yield upper limits $\eta_{i,\text{UL}}$ in Eq. (1). The upper limit $\eta_{i,\text{UL}}$ at 90\% confidence level is obtained with a Bayesian approach scanning the expected signal yield. The probability is calculated from the Gaussian-smear likelihood to take into account the systematic uncertainty.

### 4. Systematic uncertainties

The systematic uncertainties of the measurement of the $D_1(2420)^+$ resonance parameters and the Born cross sections listed in Tables 2 and 3 include correlated (common) contributions, from tracking, PID, luminosity measurements, vacuum polarization factors, interference effect and the input branching fraction, as well as uncorrelated (independent) contributions from background shapes, mass scaling, detector resolution, signal shape due to the angular distributions, and radiative corrections.

- Uncertainties of tracking and PID are each 1\% per track [28].
- The systematic uncertainties due to background contributions are estimated by leaving their magnitudes free in the fit and changing the ranges of the sideband regions. The statistical errors of the sideband samples are also included in the background uncertainty.
- The mass scale uncertainty for $D_1(2420)^+$ mass is estimated from the mass shift of $RM(D^+)$ in the control sample of $e^+e^- \rightarrow D^+D^-$. To be conservative, the largest mass shifts among the three high luminosity energy points, 0.8 MeV/$c^2$, is assigned as the systematic uncertainty due to the mass scale.
- The uncertainties due to the detector resolution are accounted for by changing the Gaussian widths for smearing the signal shape in the fit to the $RM(D^+)$ distribution. These widths, representing the resolution difference between data and MC, are varied within the uncertainty obtained from the control sample of $e^+e^- \rightarrow D^+D^-$ events. The resultant maximum changes on the numerical results are considered as the systematic uncertainties due to the detector resolution.
The systematic uncertainties are summarized in Tables 2 and 3; the sum of different uncertainties are obtained by adding up all the relevant contributions in quadrature.

5. Discussion and summary

In summary, based on $e^+e^-$ annihilation data at $E_{c.m.} = 4385.3, 4387.4, 4415.6, 4467.1, 4527.1, 4574.5$, and 4599.5 MeV, we studied the $D_1(2420)^+\pi^-$ in the mass spectrum of $D^+\pi^-\pi^-$ system in the final state of $e^+e^-\rightarrow D^+D^-\pi^+\pi^-$. The mass and width of the $D_1(2420)^+$ are measured to be $(2427.2 \pm 1.0 \pm 1.2)$ MeV/$c^2$ and $(23.2 \pm 2.3 \pm 2.3)$ MeV, respectively, which are consistent with the corresponding world-average values of $(2423.2 \pm 2.4)$ MeV/$c^2$ and $(25 \pm 6)$ MeV in PDG [13] and have better precisions. More accurate resonance parameters of the $D_1(2420)^+$ will better control the uncertainties of theoretical calculations for the $D_1(2420)\bar{D}$ and $D_1(2420)\bar{D}^*$ molecular explanations for the $Y(4260)$ and $Z_c(4430)$ states, respectively.

The Born cross sections of $e^+e^-\rightarrow D_1(2420)^+D^- + c.c. \rightarrow D^+D^-\pi^+\pi^- + e^+e^-\rightarrow \psi(3770)\pi^+\pi^- \rightarrow D^+D^-\pi^+\pi^-$ are measured as functions of the center-of-mass energy. The cross section line shape is consistent with previous BESIII measurement based on full reconstruction method [16]. There are some indications of enhanced cross sections for both processes between 4.36 and 4.42 GeV, where the reported states $Y(4360)$ and $\psi(4415)$ locate. Hence, the measured cross sections can be useful inputs to the properties of these states.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.
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