Comparison of Simulated and Measured Calcium Sparks in Intact Skeletal Muscle Fibers of the Frog

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ABSTRACT Calcium sparks in frog intact skeletal muscle fibers were modeled as stereotypical events that arise from a constant efflux of Ca\textsuperscript{2+} from a point source for a fixed period of time (e.g., 2.5 pA of Ca\textsuperscript{2+} current for 4.6 ms; 18°C). The model calculates the local changes in the concentrations of free Ca\textsuperscript{2+} and of Ca\textsuperscript{2+} bound to the major intrinsic myoplasmic Ca\textsuperscript{2+} buffers (troponin, ATP, parvalbumin, and the SR Ca\textsuperscript{2+} pump) and to the Ca\textsuperscript{2+} indicator (fluo-3). A distinctive feature of the model is the inclusion of a binding reaction between fluo-3 and myoplasmic proteins, a process that strongly affects fluo-3’s Ca\textsuperscript{2+}-reaction kinetics, its apparent diffusion constant, and hence the morphology of sparks. ΔF/F (the change in fluo-3’s fluorescence divided by its resting fluorescence) was estimated from the calculated changes in fluo-3 convolved with the microscope point-spread function. To facilitate comparisons with measured sparks, noise and other sources of variability were included in a random repetitive fashion to generate a large number of simulated sparks that could be analyzed in the same way as the measured sparks. In the initial simulations, the binding of Ca\textsuperscript{2+} to the two regulatory sites on troponin was assumed to follow identical and independent binding reactions. These simulations failed to accurately predict the falling phase of the measured sparks. A second set of simulations, which incorporated the idea of positive cooperativity in the binding of Ca\textsuperscript{2+} to troponin, produced reasonable agreement with the measurements. Under the assumption that the single channel Ca\textsuperscript{2+} current of a ryanodine receptor (RYR) is 0.5–2 pA, the results suggest that 1–5 active RYRs generate an average Ca\textsuperscript{2+} spark in a frog intact muscle fiber.

KEY WORDS: spark simulations • ryanodine receptors • fluo-3 • excitation-contraction coupling • frog muscle

INTRODUCTION

Ca\textsuperscript{2+} sparks (Cheng et al., 1993; Tsugorka et al., 1995; Klein et al., 1996) arise from brief, localized releases of Ca\textsuperscript{2+} from the SR into the myoplasm. In skeletal muscle, the release takes place at the triadic junctions through one RYR or a small cluster of RYRs. After release, Ca\textsuperscript{2+} diffuses throughout the myoplasm and binds to myoplasmic Ca\textsuperscript{2+} buffers, some of which are also able to diffuse. A quantitative model of the three-dimensional spread of Ca\textsuperscript{2+} within the myoplasm of frog muscle was first developed for Ca\textsuperscript{2+} release by an action potential (Cannell and Allen, 1984). An updated version of this model, also developed for action potential stimulation in frog fibers, provided a good description of both the spatially averaged Ca\textsuperscript{2+} signal that can be measured with low-affinity Ca\textsuperscript{2+} indicators (Baylor and Hollingworth, 1998) and the spatially resolved Ca\textsuperscript{2+} signals that can be measured with a confocal microscope and the high-affinity indicator fluo-3 (Hollingworth et al., 2000). The main goal of this article is to adapt this model for simulations of Ca\textsuperscript{2+} sparks in frog fibers and to find out whether the properties of the simulated sparks agree with those of measured sparks. Because the properties of Ca\textsuperscript{2+} sparks in frog intact fibers differ substantially from those in frog cut fibers (Hollingworth et al., 2001), the scope of this article concerns sparks in intact fibers only (18°C).

The general approach is similar to that used previously to model calcium sparks in cardiac myocytes (Prats-Sanchez and Balke, 1996; Izu et al., 1998, 2001; Smith et al., 1998) and frog cut fibers (Jiang et al., 1999; Ríos et al., 1999; Shirokova et al., 1999). Our model assumes that (a) a spark is generated by a brief, constant efflux of Ca\textsuperscript{2+} from a point source into the myoplasm; (b) the resultant increase in myoplasmic-free [Ca\textsuperscript{2+}] causes increased complexation of Ca\textsuperscript{2+} with the intrinsic Ca\textsuperscript{2+} buffers of myoplasm (ATP, parvalbumin, troponin), with the SR Ca\textsuperscript{2+} pump, and with the indicator used for the spark measurement (fluo-3); (c) gradients in the concentrations of free Ca\textsuperscript{2+} and in the Ca\textsuperscript{2+}-free and Ca\textsuperscript{2+}-bound forms of the mobile buffers (fluo-3, ATP, parvalbumin) drive diffusive movement of these species away from (or toward) the source; and (d) these processes occur isotropically in myoplasm.

The simulated calcium sparks described in this article were determined from simulated line scan images (x-t images) obtained by convolution of the microscope point-spread function (PSF) with the calculated distri-
butions of fluo-3 and Cafluo-3. Variability in the sparks was introduced by simulation of photon noise, by random selection of the time of onset of Ca$^{2+}$ release relative to the time of data sampling, and by random selection of the location of the scan line relative to that of the Ca$^{2+}$ source (see also Izu et al., 1998). For a particular set of simulated noisy sparks, the duration and amplitude of the Ca$^{2+}$ source flux were selected so that the mean values of amplitude and rise time would approximately match those of the measurements. The spatial and remaining temporal features of these sparks were then compared with the measurements to assess model performance. These features include full width at half maximum (FWHM), decay time constant, full duration at half maximum (FDHM), late baseline offset, and spark mass (a morphological parameter that emphasizes the three-dimensional nature of a spark; Sun et al., 1998; Shirokova et al., 1999; Hollingworth et al., 2001).

In the initial simulations, the two regulatory sites on troponin were assumed to bind Ca$^{2+}$ identically and independently. These simulations showed a clear discrepancy between theory and experiment: the decay of the simulated sparks had a slow phase, caused by the dissociation of Ca$^{2+}$ from troponin, that lagged the decay of measured sparks. A modified model, which incorporated positive cooperativity in the binding of Ca$^{2+}$ to the two regulatory sites on troponin (Fuchs and Bayuk, 1976; Grabarek et al., 1983), provided a satisfactory simulation of the decay phase as well as of most other properties of measured sparks. The underlying Ca$^{2+}$ release flux used in this model is 2.5 pA (units of current) for 4.6 ms. If the Ca$^{2+}$ current through a single open RYR under physiological conditions is 0.5 pA (Kettlun et al., 2000) to 2 pA (Tinker et al., 1993), the results suggest that an average Ca$^{2+}$ spark in a frog intact fiber is generated by the opening of 1–5 RYRs for 4–5 ms.

Some of the results have appeared in abstract form (Baylor et al., 2002; Chandler et al., 2002).

**Materials and Methods**

**Overview of the Spark Model**

The model is broadly similar to other models of local Ca$^{2+}$ movements near a release site inside the myoplasm (see introduction for references). Table I gives several general parameters of the model. The myoplasm is assumed to be isotropic and spherically symmetric. It is subdivided into 101 small compartments by 101 concentric spheres centered at the source of Ca$^{2+}$ entry. Ca$^{2+}$ is injected into the innermost compartment (radius, 25 nm) at a constant rate (e.g., 3 pA) for an appropriate period of time (e.g., 5 ms). The concentrations of free [Ca$^{2+}$], the myoplasmic Ca$^{2+}$ buffers (troponin, ATP, parvalbumin, the SR Ca$^{2+}$ pump), and the indicator (fluor-3) are then calculated in all compartments by the simultaneous integration of the appropriate set of first-order differential equations (Eq. A1). Table II, column 2, gives the list of abbreviations used in this paper: FDHM, full duration at half maximum; FWHM, full width at half maximum; PSF, point-spread function.

**TABLE I**

| General Parameters of the Model | 1 | 2 |
|---------------------------------|---|---|
| Temperature                     | 18°C | |
| Sarcomere length                | 3.0 μm | |
| Radius of the spark source      | 25 nm | |
| Radius of the simulation volume | 5 μm | |
| Number of simulation compartments | 101 | |
| Relative fluorescence intensities of fluo-3 | | |
| Ca$^{2+}$-free, protein-free fluo-3 | 0.005 | |
| Ca$^{2+}$-free, protein-bound fluo-3 | 0.01 | |
| Ca$^{2+}$-bound, protein-free fluo-3 | 1 | |
| Ca$^{2+}$-bound, protein-bound fluo-3 | 1 | |
| Microscope PSF                  | | |
| FWHM, and FWHM,                 | 0.2 μm | 0.3 μm |
| Spherical symmetry is assumed, with a homogeneous distribution of constituents in the resting state. The relative values of fluo-3’s fluorescence intensity are based on the in vitro measurements of Harkins et al. (1993) in the absence and presence of aldolase (the most abundant soluble muscle protein by weight). The microscope PSF is assumed to be a product of Gaussians in x, y, and z (Eq. A3). The indicated values of FWHM are similar to those measured with 0.1 μm fluorescent beads: 0.21 μm in x and y and 0.51 μm in z (Hollingworth et al., 2001).

**Reactions between Ca$^{2+}$ and its Buffers**

ATP is able to react with both Ca$^{2+}$ and Mg$^{2+}$. Since the reaction with Mg$^{2+}$ is extremely rapid, Mg$^{2+}$ and ATP can be considered to be in equilibrium, at least for the purposes of our calculations. Consequently, the reduced equivalent reaction of Ca$^{2+}$ with ATP shown in Fig. 1 A can be used (Baylor and Hollingworth, 1998). Table III A gives the apparent rate constants.

Fig. 1 B shows the multistate reaction for the binding of Ca$^{2+}$ and protein (Pr) with fluo-3 (Harkins et al., 1993; Hollingworth et al., 2000). Protein-bound fluo-3 is assumed to be immobile.
Ca²⁺ (Baylor et al., 1983), the maximal rate of Ca²⁺

which R max, the maximal turnover rate, is 19.23 s⁻¹ by the Hill equation: rate = k⁺[E] / (1 + Kᵦ[H₉2⁺])², where k⁺ is the forward rate constant and Kᵦ[H₉2⁺] is the dissociation constant for ATP.

Concentrations and Diffusion Constants Assumed in the Model

| Constituent                           | Resting concentration (µM) | Diffusion constant (10⁻⁶ cm² s⁻¹) |
|---------------------------------------|----------------------------|----------------------------------|
| Resting free [Ca²⁺]                  | 0.050                      | 3.19                             |
| Resting free [Mg²⁺]                  | 1.000                      | —                                |
| ATP                                   | 8.000                      | 1.49                             |
| Troponin (Ca²⁺ regulatory sites)      | 360                        | 0                                |
| Parvalbumin (Ca²⁺/Mg²⁺ sites)        | 1.500                      | 0.159                            |
| Protein (sites for fluo-3 binding)   | 3000                       | 0.159                            |
| SR Ca²⁺ pump (Ca²⁺ binding sites)    | 240                        | 0                                |
| Fluorescein (Fluo-3)                 | 100                        | 0                                |
| Protein-free fluo-3                  | 11.88                      | 1.51                             |
| Protein-bound fluo-3                 | 88.12                      | 0                                |

Concentrations are referred to the myoplasmic water volume (Baylor et al., 1983) and, except for Ca²⁺ and Mg²⁺, are total concentrations. At a sarcoplasmic length of 3.0 µm, the estimated troponin concentration, if referred to the water volume surrounding the thin filament, is 180 µM; with two Ca²⁺ regulatory sites per troponin molecule, the concentration of sites is 360 µM. The myoplasmic diffusion constants are for 18°C; they were adjusted from 16°C based on a Q₁₀ of 1.35. The diffusion constants for Ca²⁺, ATP, and parvalbumin are taken from Baylor and Hollingworth (1998), whereas the value for protein-free fluo-3 is based on the formula given in Zhao et al. (1996), with fluo-3's molecular weight = 765. Because free [Mg²⁺] is assumed to be constant, the value of its diffusion constant is not used. The 11.88 µM protein-free fluo-3 consists of 10.82 µM Ca²⁺-free fluo-3 and 1.06 µM Ca²⁺-bound fluo-3; the 88.12 µM protein-bound fluo-3 consists of 85.88 µM Ca²⁺-free fluo-3 and 2.24 µM Ca²⁺-bound fluo-3.

and to react with Ca²⁺ with rate constants that are substantially smaller than those assumed for protein-free fluo-3 (Table III B). The eight rate constants in Table III B satisfy k₁₁ k₁₂ k₃₂ k₃₃ = k₁₃ k₁₄ k₄₃ k₄₄, the requirement for a noncycling reaction. Fig. 1C shows the competitive reactions of Ca²⁺ and Mg²⁺ with parvalbumin (Johnson et al., 1981). The rate constants (Table III C) are taken from Baylor and Hollingworth (1998).

Fig. 1D defines a multistate reaction for the binding and transport of Ca²⁺ by the SR Ca²⁺ pump. The first three steps of this reaction are taken from the model of Inesi and de Meis (1989); they specify the sequential binding of two Ca²⁺ ions to the enzyme. The final step, the transport of Ca²⁺ into the SR, represents a lumped approximation of the remaining steps in the Inesi and de Meis (1989) model. With [Ca²⁺]ᵢ = 50 nM, the steady-state fractional occupancies of the four enzyme states in Fig. 1D are 0.9926 (E), 0.0050 (CaE), 0.0021 (CaE') and 0.0003 (Ca₂E'). The steady-state turnover rate versus free [Ca²⁺] is well described by the Hill equation: rate = R max × [Ca²⁺]³ / ([Ca²⁺]³ + Kᵦ[N]), where R max is the maximal turnover rate, is 19.23 s⁻¹, Kᵦ is 3.57 µM, and N is 1.57. With a total enzyme concentration of 120 µM (Baylor et al., 1983), the maximal rate of Ca²⁺ removal from the myoplasm is 4.6 mM Ca²⁺ ions s⁻¹ (19.23 s⁻¹ × 120 µM × 2).

Fig. 1E and F, show two reaction schemes for the binding of Ca²⁺ to the two Ca²⁺-regulatory sites on troponin. Fig. 1E shows the reaction that was used for the first spark simulations, termed Model 1. In this reaction, the two sites are identical and independent (Johnson et al., 1981; Zot and Potter, 1987); the rate constants are given in Table III E.

Fig. 1F shows the Ca²⁺-troponin reaction that was used for the final spark simulations, Model 2; the rate constants are given in Table III F. The first Ca²⁺ ion binds with low affinity (dissociation constant, Kᵦ₁ = 13 µM) and the second Ca²⁺ ion binds with high affinity (Kᵦ₂ = 0.13 µM). The steady-state fractional occupancy of the troponin sites as a function of free [Ca²⁺] is well described by the Hill equation: fractional occupancy = [Ca²⁺]² / ([Ca²⁺]² + Kᵦ[N]), in which Kᵦ = 1.30 µM, and N = 1.89. This value of Kᵦ is the geometric mean of the two basic Kᵦ₁ and Kᵦ₂ in Model 2 and is identical to the single Kᵦ of the troponin reaction in Model 1 (Table III E).

Calculation and Analysis of Noisy Sparks

The first step in the simulations was to calculate the time courses of the concentrations of calcium and the different buffer states in the 101 compartments that extend 5 µm from the spark source. [Fluo-3] was then calculated from Eq. A2 at the different distances from the source. The PSF of the microscope was approximated as a product of three Gaussian functions (Eq. A3).
with FWHM$_x$ = 0.2 µm, FWHM$_y$ = 0.2 µm, and FWHM$_z$ = 0.5 µm (Table I); x and y denote directions perpendicular to the light path and z denotes the direction parallel to the light path. The values of [FFluo] were convolved with the microscope PSF to estimate fluorescence intensity (Eq. A4). Finally, many images of noisy sparks were calculated with the help of a random number generator to mimic several components of variability in the experimental records. In the time domain, simulated spark data were extracted every 2.0 ms, with the registration of the data sampling relative to the onset of Ca$^{2+}$ release (at the spark source) varied randomly between 0.0 and 1.9 ms in multiples of 0.1 ms. In the space domain, simulated data were extracted every 0.20 µm in x, with the registration of the data sampling relative to the position of the spark source varied randomly between 0.00 and 0.15 µm in multiples of 0.05 µm. The y and z distances between the scan line and the spark source were varied randomly between 0.00 and 1.00 µm in multiples of 0.05 µm. The photon and instrument noise was simulated by the addition of random noise (standard deviation, ΔF/F = 0.26) to the x-t image.

The final x-t image, of dimensions 14 µm × 170 ms, contained a simulated noisy spark that was positioned near the spatial center of the image and started at ~50 ms. This image was analyzed in exactly the same way as an experimental x-t image (Hollingworth et al., 2001). First, the simulated image was processed by an automatic spark detection algorithm to identify putative sparks. The temporal waveform of a spark was obtained as the average of the two spatial waveforms that occurred immediately before and after the estimated time to peak of the spark. The time to peak, and the remaining spark morphological parameters, were determined from fits of standard functions to the spark waveforms in time and space (Eqs. 1 and 2 of Hollingworth et al., 2001; see also Lacampagne et al., 1999). The fitted parameters included rise time, peak amplitude, decay time constant (also called τoff), FDHM (which reflects both rise time and decay time constant), late baseline offset, FWHM, and spark mass (calculated as 1.206 × amplitude × FWHM$^2$; Hollingworth et al., 2001). Only sparks that satisfied the broad selection criteria described in Hollingworth et al. (2001) were used.

The morphological parameters of the simulated noisy sparks were compared with those of voltage-activated sparks from R. pipiens (Hollingworth et al., 2001). These sparks were measured in intact single fibers in 13 mM [K$^+$] Ringer’s (estimated membrane potential, −60 to −65 mV). As described in Hollingworth et al. (2001), raw spark amplitudes were scaled 1.3-fold to give values similar to those expected in normal Ringer’s; the factor 1.3 is the mean value of resting F in 13 mM [K$^+$] Ringer’s divided by that in normal Ringer’s.

**Detection of Sparks of Different Amplitudes as a Function of Distance from the Source**

Unlike the situation with measured sparks, with simulated sparks, the distance D between the spark source and the intersection of the scan line with the yz plane is known: D = $\sqrt{(y^2 + z^2)}$. As the value of D is increased, spark amplitude becomes smaller and the probability of detection decreases (Pratusevich and Balke, 1996; Shirokova and Rios, 1997; Izu et al., 1998; Cheng et al., 1999; Jiang et al., 1999). The relation between this probability and D was investigated in 200,000 simulations with Model 2 that used a Ca$^{2+}$ source flux of 2.5 pA for 4.6 ms (see second half of Results). Fig. 2 shows results for sparks satisfying three different acceptance criteria for the fitted value of peak amplitude: ΔF/F ≥ 0.3 (asterisks; n = 98,565); ΔF/F ≥ 0.7 (termed “larger” sparks, diamonds; n = 48,827); the largest 10% of the larger sparks (termed “in-focus” sparks, circles; n = 4,883; see last section of Results). Fig. 2 A shows the probability that a spark with the appropriate amplitude will be detected at the value of D indicated on the abscissa. For ΔF/F ≥ 0.3, D must be < 0.780 µm for the probability to exceed 0.5. For larger and in-focus sparks, the corresponding distances are 0.553 and 0.172 µm, respectively.
the larger and in-focus sparks are 0.411 and 0.152.

Fig. 2 B shows the probability that a spark with the appropriate amplitude will have a value of D ≥ the value indicated on the abscissa. The plot indicates that half of the sparks with ∆F/F ≥ 0.3 have a value of D ≤ 0.586 μm. The corresponding values of D for the larger and in-focus sparks are 0.411 and 0.132 μm, respectively.

Each amplitude criterion also has an associated "false positive" probability, i.e., a probability that a spark is recognized in an image when, in fact, there is none. When the 200,000 simulations were repeated with a Ca2⁺ source flux of zero, the corresponding numbers of sparks that satisfied the three amplitude criteria were 18, 1, and 1, respectively, instead of 98,565, 48,827, and 4,883, as indicated above. Thus, false positives do not appear to influence the results of this article, which concern larger and in-focus sparks.

Correction of an Error in an Earlier Version of the Model

Preliminary results with our spark model have been published in abstract form (Baylor et al., 2002; Chandler et al., 2002). When the manuscript for this article was submitted for publication, one of the reviewers expressed concern about a particular simulation result. This concern was addressed with some new calculations, and a programming error was discovered in the value used for the diffusion constant of Ca2⁺-free protein-free fluo-3. After correction of this error, the concern raised by the reviewer was removed. The conclusions of the abstracts were not changed, although the values of several of the reported numbers were changed slightly in the corrected calculations. For example, the values reported in the abstracts for the Ca2⁺ flux in our standard spark simulation with Model 2 were 2.7 pA for 4.4 ms compared with the values reported in this article, 2.5 pA for 4.6 ms. The corrected values of all previously published numbers affected by this error appear in columns 3 of Tables IV and V of this article.

RESULTS

This article compares measured and simulated Ca2⁺ sparks. The measured sparks were recorded in x-t scans of intact skeletal fibers from R. pipiens (Hollingworth et al., 2001; see MATERIALS AND METHODS). For the simulations, a noise-free calculation was performed under a particular set of assumptions and with particular choices for the duration and amplitude of the SR Ca2⁺ flux. From this, a large number of noisy x-t images that mimicked experimental x-t images were calculated as described in MATERIALS AND METHODS. The exact number of noisy x-t images was selected so that the final number of analyzed sparks with peak amplitude ∆F/F ≥ 0.7 (i.e., "larger" sparks) matched that of the measurements (3,176; see below). The duration and amplitude of the SR Ca2⁺ flux were adjusted in increments of 0.1 ms and 0.1 pA, respectively, so that the mean rise time and amplitude of the larger simulated sparks matched those of the measurements. The remaining morphological parameters were compared with the measured parameters to judge the overall success of the model. Although actual sparks are undoubtedly driven by SR Ca2⁺ fluxes with different values of duration and amplitude, a single duration and amplitude in the Model 2 simulations (described below) reproduced the main morphological features of the measurements surprisingly well.

Except where noted, simulations were performed with [Ca2⁺]R = 50 nM. This value is half that assumed in most other spark models, but is close to the mean value of [Ca2⁺]R estimated for the intact fibers of the spark experiments (Hollingworth et al., 2001).

Spark Simulations with Model 1

The first spark simulations used the one-step Ca2⁺-troponin binding reaction (Fig. 1 E) with the rate constants given in Table III E (i.e., the Model 1 troponin reaction). Fig. 3, A and B, show noise-free ∆F/F waveforms centered at the spark source, calculated with a Ca2⁺-release flux of 3.3 pA for 5.3 ms. These parameters were selected so that the larger simulated noisy sparks would have mean values of rise time and amplitude that approximately matched those of the measurements (4.18 ms and 1.08 ∆F/F, respectively; column 2 of Table IV). Fig. 3 A shows the temporal waveform and Fig. 3 B shows the spatial waveform at the time of peak. These traces were obtained from the simulated values of ∆[FFluo]/[FFluo]R after convolution with the microscope PSF (see
According to the autodetection program (threshold of these larger simulated sparks. Table IV, column 2, mean values spark source and the scan line (D) followed by the mean value (except for row 1) gives the corresponding values for larger measured sparks.

The corresponding mean values in Table IV, columns 2 and 3, show some similarities and some differences. The mean rise times and amplitudes are nearly identical; this is expected since the duration and amplitude of the Ca²⁺ release flux were selected to produce good agreement (preceding section). The mean value of FWHM is 14% smaller in the simulations than in the measurements (0.873 and 1.019 μm, respectively), and the mean value of spark mass (= 1.206 × amplitude × FWHM³) is 49% smaller (0.989 and 1.938 μm³, respectively). The latter difference is mainly due to the smaller average value of FWHM in the simulations in combination with the third-power dependence of mass on FWHM (see also section below on “Spark Histograms”).

The other four parameters in Table IV reflect properties of the falling phase of a spark: decay time constant, FWHM, late baseline offset, and ΔF/F at 28–48 ms. The values of these parameters are all substantially larger in the model than in the measurements. The percentage increases are: decay time constant, 74%; FWHM, 55%; late baseline offset, 378%; ΔF/F at 28–48 ms, 494%.

The difference between the simulated and measured sparks is more readily visualized with averaged traces that have reduced noise. The simulated sparks were aligned and averaged in the same way as the measured

### Table IV

Comparison of Larger Measured and Simulated Sparks (Model 1 and Modifications to Model 1)

| Parameters and units | Measurements | 50 nM [Ca²⁺]₀ | 100 nM [Ca²⁺]₀ | No Ca²⁺ pump | No ATP | Simplified fluo-3 |
|----------------------|--------------|----------------|----------------|---------------|--------|------------------|
| Distance (D) (μm)    | 0.362 ± 0.003| 0.378 ± 0.003 | 0.361 ± 0.003 | 0.329 ± 0.002 | 0.300 ± 0.002 |
| 0–100% risetime (ms) | 4.176 ± 0.034| 4.202 ± 0.029 | 4.215 ± 0.023 | 4.165 ± 0.023 | 4.148 ± 0.023 | 4.294 ± 0.024 |
| Peak amplitude (ΔF/F) | 1.082 ± 0.007| 1.087 ± 0.006 | 1.084 ± 0.006 | 1.076 ± 0.005 | 1.088 ± 0.006 | 1.089 ± 0.005 |
| Decay time constant (ms) | 4.601 ± 0.044| 7.988 ± 0.063 | 8.249 ± 0.063 | 7.568 ± 0.063 | 6.825 ± 0.058 | 10.394 ± 0.065 |
| FWHM (μm) | 5.923 ± 0.037 | 9.199 ± 0.059 | 9.540 ± 0.063 | 8.734 ± 0.057 | 8.119 ± 0.052 | 11.631 ± 0.071 |
| Late offset (ΔF/F) | 0.0125 ± 0.0010 | 0.0598 ± 0.0007 | 0.0626 ± 0.0007 | 0.0562 ± 0.0007 | 0.0606 ± 0.0007 | 0.0626 ± 0.0007 |
| FWHM (μm) | 1.019 ± 0.007 | 0.873 ± 0.004 | 0.916 ± 0.004 | 0.881 ± 0.004 | 0.799 ± 0.004 | 0.707 ± 0.003 |
| Mass (μm³) | 1.938 ± 0.047 | 0.989 ± 0.017 | 1.135 ± 0.019 | 1.016 ± 0.019 | 0.728 ± 0.019 | 0.520 ± 0.010 |
| Amplitude at 28–48 ms (ΔF/F) | 0.018 (1.03) | 0.080 (0.78) | 0.089 (0.89) | 0.079 (0.97) | 0.058 (0.79) | 0.046 (0.89) |
|                  | (0.018) | (0.017) | (0.011) | (0.009) | (0.009) | (0.132) |

Nonparenthesized entries (column 2, measured sparks; columns 3–7, simulated sparks) are the mean ± SEM values of the parameters associated with fits to 3,176 individual sparks (amplitude ΔF/F ≥ 0.7) as illustrated in Fig. 3, G–F. In row 1, D denotes the distance between the scan line and the source, which is known in the simulations (D = \(\frac{1}{3} + \frac{1}{2}\)). All simulated sparks were calculated with the Model 1 troponin reaction (Fig. 1 E) and, except for columns 4, with [Ca²⁺]₀ = 50 nM. In columns 4–7, the column heading denotes a single modification of the simulation conditions used for column 3 (see text). Parenthesized entries were obtained from fits to the temporal and spatial waveforms of averaged sparks, as illustrated in Fig. 4. All averaged sparks were formed from 1,767 individual sparks aligned at time-of-peak and center location (see Hollingworth et al., 2001); amplitude at 28–48 ms after peak was obtained from Gaussian fits of the type illustrated in Fig. 4, C and F.
sparks, based on alignment of the estimated time of peak and the estimated spatial center of the sparks (Fig. 10 of Hollingworth et al., 2001). Fig. 4 shows temporal and spatial waveforms obtained from these averaged sparks (Fig. 4, A–C, simulations, asterisks; Fig. 4, D–F, measurements, squares). The curves in Fig. 4 are fits of the standard functions. The parenthesized values in Table IV, columns 2 and 3, give morphological parameters obtained from these fits (see also values printed on the figure panels).

In the time domain (Fig. 4, A and D), the simulated points clearly decay more slowly than the measured data; in addition, the standard function provides a poor fit in A but a good fit in D. In the spatial domain at the time of peak of the spark (Fig. 4, B and E), both the simulated points and the data are reasonably well described by gaussian functions (curves). The fitted value of FWHM is somewhat smaller in the simulations than in the measurements (0.84 and 0.96 μm, respectively). In the spatial domain 28–48 ms after time of peak (Fig. 4, C and F), the amplitude of the simulated points is about sixfold larger than that of the data (0.107 and 0.018, respectively; last row in Table IV). The value of FWHM is somewhat smaller in the simulations than in the measurements (2.11 and 2.94 μm, respectively; see figure panels).

The substantial differences between the values of decay time constant, FDHM, late baseline offset, and ΔF/F at 28–48 ms in Table IV, columns 2 and 3, indicate that one or more of the processes that determine the falling phase of a spark are not well described by Model 1. It was therefore of interest to find out whether a simple modification of the model might provide better agreement with the measurements.

**Modifications of Model 1**

Model 1 differs in several ways from most previously published spark models. The differences include: (a)
the value assumed for $[\text{Ca}^{2+}]_R$; (b) the particular reaction chosen for the SR $\text{Ca}^{2+}$ pump; (c) the inclusion of ATP as an endogenous $\text{Ca}^{2+}$ buffer; and (d) the inclusion of a binding reaction between fluo-3 and myoplasmic protein. To evaluate the effect of these differences, the parameters of the model were modified as described in the following paragraphs. In each case, the amplitude of the $\text{Ca}^{2+}$ flux at the source was readjusted so that the mean amplitude of the larger simulated noisy sparks matched that of the measurements. Table IV, columns 4–7 give the morphological parameters obtained in the simulations after these modifications.

The first modification was to increase the value of $[\text{Ca}^{2+}]_R$ from 50 to 100 nM. The main effect of this change is to increase resting fluorescence ($F$) by the factor 1.73. The fractional increase in $F$ is smaller than that in $[\text{Ca}^{2+}]_R$ because the fluorescence of $\text{Ca}^{2+}$-free fluo-3 is not negligible and the binding of $\text{Ca}^{2+}$ to fluo-3 varies slightly nonlinearly with $[\text{Ca}^{2+}]_R$. To maintain good agreement between the mean amplitude of the simulated noisy sparks and that of the measurements, it was necessary to increase the amplitude of the SR $\text{Ca}^{2+}$ flux from 3.3 to 6.0 pA. With this change, the mean values of all remaining morphological parameters are very similar to those with $[\text{Ca}^{2+}]_R = 50$ nM (Table IV, columns 3 and 4). The largest changes are a 5% increase in FWHM and a 15% increase in spark mass (which has a third-power dependence on FWHM). An increase in FWHM with an increase in $[\text{Ca}^{2+}]_R$ is expected, because an increase in $[\text{Ca}^{2+}]_R$ should reduce the resting concentration of $\text{Ca}^{2+}$-free binding sites on troponin and the SR $\text{Ca}^{2+}$ pump (which are immobile $\text{Ca}^{2+}$ buffers) as well as on parvalbumin (which is only slightly mobile). This reduction of $\text{Ca}^{2+}$-free binding sites on molecules with restricted mobility, and the twofold increase in $\text{Ca}^{2+}$ source flux, should allow free $\text{Ca}^{2+}$, $\text{CaATP}$, $\text{Cafluo-3}$, and hence fluo-3 fluorescence to spread more rapidly from the source. Overall, the good agreement between the values in columns 3 and 4 indicates that the differences between Model 1 (column 3) and the measurements (column 2) are not resolved by an increase in $[\text{Ca}^{2+}]_R$ from 50 to 100 nM.

The next modification was to remove the SR $\text{Ca}^{2+}$ pump (with $[\text{Ca}^{2+}]_R$ restored to 50 nM). This change required a 9% reduction in $\text{Ca}^{2+}$ release to maintain the mean amplitude of the larger simulated sparks at $\Delta F/F = 1.08$. The values of the other morphological parameters in Table IV, column 5, are similar to those in column 3. The largest changes are 5–6% reductions in decay time constant and late baseline offset. The qualitative explanation is that, with the pump present, some of the $\text{Ca}^{2+}$ that associates with the pump during
the rising and early falling phase of the spark is able to dissociate from the pump during the later falling phase (see steps 1–3 in Fig. 1 D). This dissociation of Ca\(^{2+}\) elevates free [Ca\(^{2+}\)] and hence fluo-3 fluorescence, thereby prolonging the spark. Overall, the relatively minor changes in morphological parameters associated with removal of the Ca\(^{2+}\) pump make it unlikely that a different reaction scheme for the pump would remove the differences between the values in columns 2 and 3 of Table IV.

The third modification was to set the ATP concentration to zero. Similar to the removal of the SR Ca\(^{2+}\) pump, the removal of ATP is expected to decrease the Ca\(^{2+}\) flux required to maintain a particular spark amplitude and to decrease the decay time constant of the spark. In addition, because ATP constitutes a substantial pool of rapidly reacting, mobile Ca\(^{2+}\) buffer (Baylor and Hollingworth, 1998), removal of ATP is expected to decrease spark FWHM. As shown in Table IV, columns 4 and 6, these expectations are realized. With ATP removed, the Ca\(^{2+}\) flux is reduced by 33% and the mean values of decay time constant and FWHM are reduced by 15% and 11%, respectively; the reduction in spark mass is 26%. These differences in morphological parameters do not reconcile the differences between Model 1 and the measurements (Table IV, columns 2 and 3).

The fourth modification was the elimination of protein from the fluo-3 reactions in Fig. 1 B (“simplified fluo-3” simulation; Table IV, column 7). With only the one-step reaction between Ca\(^{2+}\) and fluo-3 remaining, three other changes were required to make the properties of fluo-3 in the model consistent with those in previous measurements. \(k_{-1}\) and \(k_{-1}\) (Fig. 1 B) were set to \(0.402 \times 10^8\) M\(^{-1}\) s\(^{-1}\) and 63.2 s\(^{-1}\), respectively, and the diffusion constant of fluo-3 was set to \(0.212 \times 10^{-6}\) cm\(^2\) s\(^{-1}\). These are the apparent values estimated to apply to single effective pools of fluo-3 and Cafluo-3 in intact fibers at 18\(^\circ\)C, based on the indicator’s apparent diffusion constant measured under resting conditions and its apparent reaction rate constants with Ca\(^{2+}\) measured under action potential conditions (Harkins et al., 1993; Baylor and Hollingworth, 1998). With these modifications, a Ca\(^{2+}\) flux of 1.7 pA (rather than 3.3 pA) is required to maintain the mean spark amplitude near the measured value. The other morphological parameters, however, are changed so as to exaggerate the differences between the simulations and the measurements. Thus, in agreement with the conclusion of Hollingworth et al. (2000), the behavior of fluo-3 in myoplasm is better described by a model that includes reactions between fluo-3 and protein (Fig. 1 B) than by a model that has only single effective pools of fluo-3 and Cafluo-3.

Overall, none of the four modifications in Table IV substantially reduced the discrepancy between the falling phases of the simulated and measured sparks. Furthermore, there is no indication that any combination of the four modifications would reduce the discrepancy. Our conclusion is that some other feature of the model must be modified.

**Spark Model 2**

At late times, the decay of the sparks simulated with Model 1 is slower than that of the measurements. The likely explanation is that, in the simulations, the Ca\(^{2+}\) that is captured by troponin during the rising and early falling phase of the spark is released during the later falling phase. This released Ca\(^{2+}\) elevates free [Ca\(^{2+}\)] and prolongs the decay of [CaFluo] and [CaPrFluo] (Fig. 1 B). Better agreement between the simulations and the measurements might be achieved with a modified troponin reaction in which troponin binds Ca\(^{2+}\) with a greater delay during the rising and early falling phase of the spark and/or releases it more slowly during the later falling phase.

The troponin reaction used in spark Model 2 (Fig. 1 F with the reaction rate constants in Table III F) has these kinetic features. The first Ca\(^{2+}\) ion binds with low affinity and the second ion binds with high affinity. There is a substantial kinetic delay between an increase in [Ca\(^{2+}\)] and an increase in the concentration of Ca\(^{2+}\) bound to the high-affinity, slowly dissociating site. This reduces the amount of Ca\(^{2+}\) captured by troponin during the rising and early falling phase of a spark. Then, the slow dissociation of Ca\(^{2+}\) from the high-affinity site decreases the rate at which Ca\(^{2+}\) is released from troponin during the late falling phase of a spark. Hence, [Ca\(^{2+}\)], [CaFluo], and [CaPrFluo] are expected to decline more rapidly in Model 2 than in Model 1.

Another difference between the reactions in Fig. 1, E and F, concerns steady-state cooperativity. Ca\(^{2+}\) binding proteins like troponin that have two or more Ca\(^{2+}\) binding sites of the “E-F hand” configuration, often bind Ca\(^{2+}\) with some degree of positive cooperativity (e.g., Ames et al., 1995). Indeed, positive cooperativity has been detected in the binding of Ca\(^{2+}\) to the myofilaments of skeletal muscle (rabbit psoas muscle, pCa’s in the range 7 to 5; Fuchs and Bayuk, 1976) and to the regulatory sites of fluorescently labeled troponin reconstituted onto thin filaments (rabbit back and leg muscle; Grabarek et al., 1983). The Ca\(^{2+}\) binding reaction in Fig. 1 F also shows positive cooperativity (steady-state Hill coefficient, 1.89; see MATERIALS AND METHODS), whereas the reaction in Fig. 1 E is noncooperative.

The use of the Model 2 troponin reaction is clearly speculative. The reaction between Ca\(^{2+}\) and the regulatory sites on troponin is quite sensitive to experimental conditions (e.g., Grabarek et al., 1983; Rosenfeld and Taylor, 1985; Zot et al., 1986) and the reaction that applies under physiological conditions remains to be determined. Nonetheless, the qualitative expectations de-
scribed above for the Model 2 troponin reaction encouraged us to explore its properties more fully.

Fig. 5, A and B, illustrate some of the kinetic differences between the Model 1 and Model 2 troponin reactions. In Fig. 5 A, the top trace shows $[\text{Ca}^{2+}]$, which is assumed to change in a spatially uniform manner (non-spark situation) from a resting level of 50 nM to 5 $\mu$M for a period of 5 ms. The superimposed pair of traces in the middle shows $f$, the fractional occupancy of the troponin regulatory sites with $\text{Ca}^{2+}$ (continuous trace, Model 1 reaction; dotted trace, Model 2 reaction). The value of $f$ in the Model 2 reaction has a smaller resting value than that in the Model 1 reaction (0.003 vs. 0.037), a slower rising phase in response to the increase in $\text{Ca}^{2+}$, a smaller peak value (0.515 vs. 0.763), and a markedly slower decline in response to the decrease in $\text{Ca}^{2+}$. The superimposed pair of traces at the bottom shows $df/dt$. After the pulse of $\text{Ca}^{2+}$, $df/dt$ in the Model 1 reaction shows a substantial negative phase that lasts $\sim$10 ms, whereas, in the Model 2 reaction, the late phase is greatly reduced. Thus, a Model 2 spark would be expected to have a briefer falling phase than a Model 1 spark.

In Fig. 5 B, the upper trace shows the spatially averaged $[\text{Ca}^{2+}]$ waveform estimated for a frog intact fiber stimulated by a single action potential (Baylor and Hollingworth, 1998). The middle traces ($f$) show that the troponin sites reach nearly full saturation with both the Model 1 and 2 reactions in response to the normal physiological stimulus (peak fractional occupancies, 0.922 and 0.961, respectively). The main difference occurs during the falling phase, where the time course of $\text{Ca}^{2+}$’s dissociation from troponin is noticeably slower in the Model 2 reaction.

**Spark Simulations with Model 2**

Table V and Figs. 6 and 7 give results for spark simulations performed with Model 2. These are presented in the same formats used for Model 1 (Table IV, Figs. 3, A–D, and 4). As usual, the duration and amplitude of the SR $\text{Ca}^{2+}$ flux were adjusted so that the mean rise time and amplitude of the larger simulated sparks matched those of the measurements. The flux values used for the standard Model 2 simulation (2.5 pA and 4.6 ms; Table V, column 3) are both smaller than the corresponding values used for the standard Model 1 simulations (3.3 pA and 5.3 ms). Compared with Model 1, the Model 2 values of the mean morphological parameters are in substantially better agreement with the measurements (Table V, columns 2 and 3), including the values of FWHM, mass, and the problematic parameters that reflect the falling phase of the spark (decay time constant, FDHM, late baseline offset, $\Delta F/F$ at 28–48 ms).
The improvement with Model 2 can also be seen in a comparison of the temporal and spatial waveforms of the averaged sparks (Fig. 7). Although the temporal waveform calculated with Model 1 is poorly fitted by the standard function (curve in Fig. 4 A), the waveform calculated with Model 2 is well fitted (curve in Fig. 7 A), as is the measured waveform (curve in Fig. 7 D). The values of the fitted parameters in Fig. 7, A and D, are also in good agreement. The main discrepancy between the points and the fitted curves in Fig. 7, A and D, occurs 10–40 ms after the peak, where the points lie slightly above the curves. All three spatial waveforms at time of peak are well-fitted by the standard function, although the value of FWHM obtained in the Model 1 fit (0.84 \( \mu m \); Fig. 4 B) is noticeably smaller than those obtained in the Model 2 fit (0.97 \( \mu m \); Fig. 7 B) and the

The format is identical to that of Table IV, as are the values in column 2. Columns 3–7 give simulated data obtained with Model 2, i.e., the two-step tropo-

### Table V
Comparison of Larger Measured and Simulated Sparks (Model 2 and Modifications to Model 2)

| Parameters and units | Measurements | 50 nM [Ca\(^{2+}\)]\(_R\) | 100 nM [Ca\(^{2+}\)]\(_R\) | No Ca\(^{2+}\) pump | No ATP | Simplified fluo-3 |
|----------------------|--------------|----------------------|----------------------|---------------------|--------|------------------|
| Distance (D) (\( \mu m \)) | 0.405 ± 0.003 | 0.420 ± 0.003 | 0.430 ± 0.003 | 0.348 ± 0.002 | 0.350 ± 0.002 |
| 0–100% risetime (ms) | 4.176 ± 0.034 | 4.194 ± 0.023 | 4.159 ± 0.023 | 4.175 ± 0.023 | 4.020 ± 0.021 | 4.389 ± 0.026 |
| Peak amplitude (\( \Delta F/F \)) | 1.082 ± 0.007 | 1.077 ± 0.005 | 1.081 ± 0.006 | 1.083 ± 0.006 | 1.079 ± 0.005 | 1.083 ± 0.005 |
| Decay time constant (ms) | 4.601 ± 0.044 | 5.747 ± 0.037 | 5.981 ± 0.040 | 4.706 ± 0.033 | 4.933 ± 0.034 | 7.270 ± 0.042 |
| FDHM (ms) | 5.923 ± 0.037 | 6.820 ± 0.030 | 7.044 ± 0.032 | 6.012 ± 0.025 | 6.103 ± 0.026 | 8.097 ± 0.035 |
| Late offset (\( \Delta F/F \)) | 0.0125 ± 0.0010 | 0.0244 ± 0.0006 | 0.0314 ± 0.0006 | 0.0202 ± 0.0006 | 0.0275 ± 0.0006 | 0.0243 ± 0.0007 |
| FWHM (\( \mu m \)) | 1.019 ± 0.007 | 1.014 ± 0.005 | 1.040 ± 0.005 | 1.067 ± 0.005 | 0.855 ± 0.004 | 0.852 ± 0.004 |
| Mass (\( \mu m^3 \)) | 1.958 ± 0.047 | 1.528 ± 0.027 | 1.646 ± 0.026 | 1.784 ± 0.028 | 0.947 ± 0.019 | 0.922 ± 0.019 |
| Amplitude at 28–48 ms (\( \Delta F/F \)) | (0.018) | (0.030) | (0.040) | (0.023) | (0.031) | (0.041) |

The format is identical to that of Table IV, as are the values in column 2. Columns 3–7 give simulated data obtained with Model 2, i.e., the two-step tropo-

**Figure 6.** Simulated spark waveforms calculated with Model 2 and a Sr Ca\(^{2+}\) flux of 2.5 pA for 4.6 ms; the format is the same as that in Fig. 3, A–D. In C and D, the distance between the scan line and the spark source is 0.20 \( \mu m \) in y and 0.35 \( \mu m \) in z. In C, the value of FDHM is 6.69 ms.
measurements (0.96 µm; Fig. 7 E). In all three of these spatial fits, the points near 0 µm and near ±1 µm lie slightly above the curves. Improvement with Model 2 is also apparent in the amplitude of the spatial waveform 28–48 ms after the spark peak. The amplitude of the simulated waveform (0.030, last row in Table V; Fig. 7 C) is substantially closer to that of the measured waveform (0.018; Fig. 7 F) than is the amplitude obtained with Model 1 (0.107; Fig. 4 C).

The four modifications that were applied to Model 1 (Table IV, columns 4–7) gave similar results when applied to Model 2 (Table V, columns 4–7). In Model 2, the reduction in decay time constant associated with removal of the Ca$^{2+}$ pump is fractionally somewhat greater than that in Model 1 (18% and 5%, respectively). This occurs because, in Model 2, troponin releases Ca$^{2+}$ at a smaller rate during the falling phase of the spark, so that the dissociation of Ca$^{2+}$ from the other buffers has a more pronounced effect.

A clear limitation of our simulations is that they use a stereotypical SR Ca$^{2+}$ release flux (fixed duration and amplitude), whereas the flux that underlies actual sparks undoubtedly involves a range of durations and amplitudes. In addition, some of the reactions between Ca$^{2+}$ and its buffers, and between fluo-3 and protein, may be more complicated than those shown in Fig. 1. For example, the troponin reaction in Fig. 1 F is the only alternate troponin reaction explored in our simulations and it would be surprising if future studies did not show that the actual reaction is more complicated. On balance, it is remarkable that Model 2, with its many simplifications, is able to reproduce the average properties of measured sparks as well as it does.

**Histograms of Spark Morphological Parameters**

In Table V, the values of SEM for the simulations are smaller than those for the measurements, indicating less dispersion in the simulated parameters. Fig. 8 gives additional information about this difference in dispersions. In each panel, the normalized histograms show the relative frequency of occurrence of a particular value of a morphological parameter (simulated sparks, asterisks; measured sparks, open squares).

Fig. 8 A shows the rise time histograms. The value of rise time is thought to closely reflect the duration of the underlying SR calcium release flux responsible for a spark (Cannell et al., 1995). For all simulated sparks, the flux duration was exactly 4.6 ms; thus, the dispersion in the simulated histogram simply reflects the combined effects of the various sources of variability and noise that were used to generate the simulated noisy sparks. The simulated histogram has fewer events at <3 ms and at >6 ms than the measured histogram.
and a concentration of events between 3 and 6 ms, near the flux duration (4.6 ms). There are two reasons why the actual difference at <3 ms may be more pronounced than that shown in Fig. 8 A. First, the measured histogram may be biased against the inclusion of events that arise from Ca²⁺ source fluxes of short duration. Such events are likely to be of smaller amplitude (Lacampagne et al., 2000; see also next section) and consequently detected less reliably by our autodetection program. Second, with 2 ms sampling, the rise time of events with a source flux duration <3 ms is expected to be overestimated (Hollingworth et al., 2001). Because of these biases, and the noise level associated with the data points in the measured histogram, no attempt was made to express the measured histogram in terms of a weighted sum of simulated histograms generated with different flux durations. The general similarity of the histograms in Fig. 8 A, however, suggests that the duration of SR Ca²⁺ release for most larger sparks measured under the conditions of our experiments is narrowly distributed with a mode of 4–5 ms.

Fig. 8 B shows the amplitude histograms. The simulated and measured points are in reasonable agreement, as expected, because of the strong dependence of amplitude on the distance D between the scan line and the source. This dependence should be similar for both simulations and measurements (see also “Properties of In-focus Sparks”).

Fig. 8 C shows the histograms of tau_{off} (decay time constant). The simulated histogram is right-shifted with respect to the measured histogram, as expected from the mean values in Table V. It also shows slightly less dispersion, which may be due to both intrafiber and interfiber variation in the processes that determine tau_{off} in the measurements.

Fig. 8 D shows the histograms of late offset. Both histograms have a mode at ΔF/F = 0.02, as expected from the mean values in Table V. The greater dispersion in the measured histogram may reflect variability in fiber properties. It may also reflect small elevations in [Ca²⁺] produced by subthreshold sparks in the same sarcomere as a detected spark. Such increases in [Ca²⁺] could bias...
the estimate of the late offset either up or down, depending on whether the subthreshold activity followed or preceded, respectively, the detected spark.

Fig. 8 E shows the histograms of FWHM. Both histograms have modes near 1.0 μm, but the simulated histogram has less dispersion. The extra events in the measured histogram at FWHM values <0.7 μm are particularly noticeable. They may arise, at least in part, from events with rise times <3 ms (Fig. 8A); these allow less time for the diffusive spread of Ca2+, fluo-3, and the other mobile Ca2+ buffers from the source at the time FWHM is estimated.

Fig. 8 F shows the histograms of spark mass. The simulated histogram has a well-defined mode near 1 μm³, whereas the measured histogram shows a significant fraction of events with mass <0.5 μm³. These events undoubtedly arise from measured sparks with small values of FWHM (Fig. 8 E).

The general conclusion of this section is that the value of each morphological parameter obtained from a collection of either simulated or measured sparks shows a large degree of dispersion. With simulated sparks, this occurs even though the model parameters and source flux are constant. In this situation, the dispersion must arise from the simulation of random photon and instrumentation noise, from the random selection of the time of onset of Ca2+ release relative to the time of data sampling, and from the random selection of the location of the scan line relative to that of the Ca2+ source. As shown in Fig. 8, the dispersion of parameters from measured sparks is somewhat greater than that from simulated sparks. Although the factors that are responsible for the dispersion in the simulations are expected to make similar contributions to the dispersion in the measurements, an additional component(s), due to variability in other factors, appears to be present in the measurements. Such factors include the following: (a) the duration and amplitude of the source flux of Ca2+, (b) [Ca2+]R, (c) the concentrations of the myoplasmic Ca2+ buffers (troponin, parvalbumin, the Ca2+ pump, ATP, and fluo-3), (d) anatomic structures that may affect spark spread (the triadic junctions, the myofibrils, the SR, mitochondria, etc.); (e) small fluctuations in [Ca2+] caused by subthreshold sparks.

**Effect of Variations In Flux Duration and Amplitude on Spark Morphological Parameters**

The extra dispersion in the measured rise-time histogram at <3 ms and >6 ms (Fig. 8A) suggests that actual sparks have variable durations of Ca2+ source flux. To explore the effect of source duration on the fitted morphological parameters, simulations were done with Model 2 with the standard flux amplitude of 2.5 pA (Table V, column 3) and three additional flux durations (2.0, 3.0 and 6.0 ms) (Table VI, columns 2–4). As the flux duration is increased, the amplitude of ΔF/F is increased, and the likelihood of detection of sparks whose source is away from the scan line is increased. Consequently, the mean value of D is increased (Table VI, row 1). Larger flux durations also increase the mean values of all other morphological parameters. For a twofold increase in duration (from 3.0 to 6.0 ms; columns 3 and 4), the percentage increases in mean parameter values are as follows: rise time (72%), peak amplitude (14%), decay time constant (35%), FDHM (56%), late offset (16%), FWHM (29%), and spark mass (106%).
The close correlation between flux duration and rise time is expected from the simulations of Cannell et al. (1995). Since rise time, but not flux duration, can be directly estimated from fits to measured sparks, plots were made of various morphological parameters versus rise time. Unfortunately, the plots were noisy and, for each parameter, were not very different for the measurements and the standard simulations with Model 2. With a larger population of measured sparks and/or with more rapid data sampling, correlations between rise time and other morphological parameters of the type suggested in Table VI, columns 2–4, might become detectable.

To examine the effect of flux amplitude on spark properties, simulations were performed with Model 2 at the standard flux duration of 4.6 ms and three additional flux amplitudes (2.0, 4.0, and 5.0 pA) (Table VI, columns 5–7). An increase in flux amplitude, similar to an increase in flux duration, increases spark amplitude. This again increases the likelihood of detection of sparks whose source is away from the scan line, so that the mean value of D is increased. The mean values of the other morphological parameters also increase with flux amplitude. For a twofold increase in flux amplitude (2.0 to 4.0 pA), the increases are as follows: rise time (6%), peak amplitude (30%), decay time constant (28%), FDHM (21%), late offset (46%), FWHM (15%), and spark mass (69%). The 30% increase in spark amplitude is substantially smaller than the 100% increase in flux amplitude.

**Properties of In-focus Sparks**

As noted in materials and methods, the peak amplitude of a spark decreases as D increases. Other morphological parameters may also depend on D. To reduce such variability, sparks that arose near the Ca$^{2+}$ source were investigated. These were tentatively identified as the top 10% of the larger sparks (referred to as “in-focus” sparks). In the standard simulations with Model 2, the average value of D for in-focus sparks is 0.15 μm, which is similar to the spacing of the pixels in the measurements, 0.20 μm; in contrast, with larger sparks, the average value of D is 0.41 μm (Table V, column 3).

Table VII compares the properties of in-focus simulated and measured sparks. As expected, the mean peak amplitudes of the in-focus sparks (1.878 and 1.703, Table VII, columns 2 and 3) are substantially larger than those of the larger sparks (1.082 and 1.077, Table V, columns 2 and 3). Overall, the mean morphological values of the in-focus simulated sparks are in fair agreement with those of the measurements. The largest discrepancy concerns spark mass, which is 1.272 μm$^3$ in the simulations and 2.658 μm$^3$ in the measurements. This difference arises because in-focus simulated sparks have, on average, both a smaller peak amplitude and a smaller FWHM than in-focus measured sparks. Although the reason for these latter differences is not known, it may be of interest to explore the possibility that they arise from the expected variability of Ca$^{2+}$ source fluxes in the measured, but not the simulated, sparks.

Fig. 9 shows temporal and spatial waveforms for averaged in-focus sparks (simulations, asterisks; measurements, open squares). All waveforms are well-fitted by the standard functions (curves). The morphological parameters obtained from the fits (see figure panels and parenthetical entries in Table VII) indicate that the simulations and the measurements are in reasonable agreement.

The general agreement between columns 2 and 3 in Table VII, and between the averaged waveforms in Fig. 9, A–C and D–F, adds support to the conclusion that Model 2 provides a reasonable simulation of the measurements.

**Discussion**

This article describes a model of Ca$^{2+}$ sparks in frog intact skeletal muscle fibers and compares the morphological properties of simulated and measured sparks. Many of the basic features of our model are similar to those of previously published models of sparks and other local Ca$^{2+}$ movements in skeletal and cardiac muscle (see introduction). Consequently, the properties of our calculated sparks share many similarities.
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with those calculated with the other models, such as a value of FWHM of \~1.0 \(\mu\)m. Because the average value of FWHM measured in frog cut fibers and cardiac myocytes is 1.3–2.5 \(\mu\)m (see Discussion and Table VII of Hollingworth et al., 2001), it is clear that the models do not reproduce the FWHM values of those measurements. A different situation arises, however, when our Model 2 simulations are compared with our intact fiber measurements: for the first time, reasonable agreement is found between the morphological properties of simulated and measured sparks, including the value of FWHM (columns two and three of Tables V and VII).

Our simulations incorporate several novel features not included in previous spark models in skeletal muscle. First, our simulations take into account the effects of random variation on the estimation of the morphological parameters of sparks. These include photon and instrumentation noise as well as variability in spark onset and spark location. This approach permits a more realistic comparison of simulated sparks with measured sparks. Second, our model includes binding reactions between fluo-3 and myoplasmic proteins (Fig. 1 B; see also Izu et al., 2001). In a previous study of spatially resolved fluo-3 Ca\textsuperscript{2+} signals in frog muscle (Hollingworth et al., 2000), model calculations that included these reactions produced better agreement with the fluorescence measurements than calculations that assumed single pools of fluo-3 and Cafluo-3. As expected from that study, the inclusion of the fluo-3-protein reactions in our spark model produced better agreement with the measurements than the simplified two-state fluo-3 reaction (compare columns 2, 3, and 7 of Tables IV and V). Better agreement with measured sparks after inclusion of the reaction between fluo-3 and protein was also reported in a recent model of Ca\textsuperscript{2+} sparks in cardiac myocytes (Soeller and Cannell, 2002). Third, our model includes a multistate reaction between Ca\textsuperscript{2+} and troponin (next section).

\textbf{Ca\textsuperscript{2+} Troponin Binding}

The properties of calcium sparks are determined by the duration and amplitude of the SR Ca\textsuperscript{2+} release flux, by the binding reactions between Ca\textsuperscript{2+} and the myoplasmic Ca\textsuperscript{2+} buffers (as well as Ca\textsuperscript{2+} transport by the SR Ca\textsuperscript{2+} pump), and by the diffusion of Ca\textsuperscript{2+} and the mobile Ca\textsuperscript{2+} buffers. The results of this article suggest that the decay phase of a spark in a frog intact muscle fiber (average sarcomere length, 3.0 \(\mu\)m) contains useful information about the binding of Ca\textsuperscript{2+} to troponin. In the initial simulations (Model 1), the reactions between Ca\textsuperscript{2+} and the two regulatory sites were assumed to be identical and independent (Fig. 1 E). With this assumption, the morphological parameters that reflect the spark falling phase (decay time constant, FDHM, late
baseline offset, and ΔF/F at 28–48 ms) were substantially larger in the simulations than in the measurements (Table IV, columns 2 and 3). In the Model 2 simulations, a two-step binding reaction between Ca\(^{2+}\) and troponin was used (Fig. 1 F, with rate constants given in Table III F). With Model 2, the values of decay time constant, FDHM, ΔF/F at 28–48 ms, and late baseline offset came into good agreement with the measurements, as did the value of FWHM (Table V, columns 2 and 3).

Based on the substantial differences between the Model 1 simulations and the measurements, and on the substantial reduction in these differences with Model 2, we speculate that the binding of Ca\(^{2+}\) to troponin in frog intact skeletal muscle is poorly described by the one-step independent reaction (Fig. 1 E). This possibility has not been deduced from previous Ca\(^{2+}\) indicator studies in skeletal muscle. The alternative proposal considered here, that the Ca\(^{2+}\)-troponin reaction is a two-step reaction with positive cooperativity, is supported by some work in the literature (Fuchs and Bayuk, 1976; Grabarek et al., 1983) and merits further investigation.

**Calcium Source Flux and the Number of Active RYRs Required Per Spark**

This article shows that there is good agreement between Model 2 simulated sparks and experimental sparks (Tables V and VII and Figs. 7–9). This suggests that the Ca\(^{2+}\) source flux that was used for the simulations, 2.5 pA for 4.6 ms (18°C), is similar to that responsible for an average voltage-activated spark in a frog intact muscle fiber under our measurement conditions. At physiological ion concentrations, the single-channel RYR Ca\(^{2+}\) current estimated from bilayer experiments varies between 0.5 pA (Kettlun et al., 2000) and 2 pA (Tinker et al., 1993). If the channel open time during a spark is not interrupted by flickering, these single channel currents indicate that 1–5 RYRs would need to open to give 2.5 pA of Ca\(^{2+}\) source flux for a spark in a frog intact skeletal fiber.

Estimates of the number of RYRs per spark have also been made in frog cut fibers, where the properties of sparks differ significantly from those in intact fibers (see Table VII of Hollingworth et al., 2001). The numbers estimated for cut-fiber sparks cover a wide range: ≤4 (Shifman et al., 2000), ≥6 (Gonzalez et al., 2000, and 8–60 (Rios et al., 1999). Our estimate for intact fibers is close to that of Shifman et al. (2000).

**Appendix**

**Numerical Simulation of the Concentrations of Calcium and Calcium Buffers Near a Point Source of Calcium Flux from the SR**

The theoretical approach for these simulations is similar to that used by Cannell and Allen (1984), Pratusevich and Balke (1996), Baylor and Hollingworth (1998), Izu et al. (1998, 2001), Smith et al. (1998), Jiang et al. (1999), Rios et al. (1999), and Shirokova et al. (1999). Ca\(^{2+}\) efflux from a small, almost point source is assumed to start at zero time, to be of constant amplitude, and to stop after a few milliseconds. After Ca\(^{2+}\) exits from the source, it diffuses in the myoplasmic solution and reacts with several different calcium buffers according to the reactions illustrated in Fig. 1: ATP (A), the calcium indicator fluo-3 (B), parvalbumin (C), the SR Ca\(^{2+}\) pump (D), and troponin (E or F). Table II gives the resting concentrations of the buffers and their ligands, which are assumed to be constant in space. Table III gives the values of the rate constants of the various reaction steps. The medium surrounding the source is taken to be isotropic and isopotential. Under these conditions, the equations that describe the changes in concentrations of Ca\(^{2+}\) and the buffer states can each be written in the form,

$$\frac{\partial c}{\partial t} = D \left( \frac{\partial^2 c}{\partial r^2} + \frac{2 \partial c}{\partial r} \right) + R_{cr}, \quad (A1)$$

in which c is the concentration of Ca\(^{2+}\) or one of the buffer states, t is time, D is the diffusion constant (equal to 0 if the substance is immobile), r is the distance from the center of the calcium source, and R\(_{cr}\) is the rate at which c is generated by the reactions illustrated in Fig. 1. In the case of Ca\(^{2+}\) itself, R\(_{cr}\) also includes any Ca\(^{2+}\) that exits from the source. The first term on the right-hand side of Eq. A1 gives the well known rate of change for radial diffusion (Crank, 1956).

These equations for the various c\(_i\)s were solved numerically. The volume surrounding the calcium source out to 5.025 μm was separated into 101 regions by the same number of concentric spherical shells. The radius of the first shell, r\(_0\)(0), is 0.025 μm and the separation between shells is 0.05 μm. Accordingly, r\(_i\)(0) = (0.025 + 0.05i) μm, i = 0, 1,..., 100. The volume of the space inside the spherical shell at r\(_i\)(0), denoted by v\(_i\)(0), is equal to (4/3)πr\(^3\)\(_i\)(0). The volume of the other spaces, denoted by v\(_i\)(i), is given by

$$v(i) = \frac{4}{3} \pi [r_i^3(i) - r_i^3(i-1)], \quad i = 1, 2, ..., 100.$$

Within each volume v\(_i\)(i), the concentration of each substance c is assumed to be constant with a value denoted by c\(_i\)(i) and a position v\(_i\)(i) equal to 0.05i μm, i = 0, 1,..., 100. The source flux of Ca\(^{2+}\) is assumed to enter v\(_0\)(0). The other boundary condition requires that all concentrations inside v\(_{100}\) are equal to their respective resting values.

A series of differential equations in time were used to describe the changes in the various c\(_i\)(i)s for i = 0, 1,..., 99. These were solved by numerical integration with a modified Runge-Kutta routine (Romanelli, 1960) with an integration step of 1 μs. The total integration time was typically 120 ms after the beginning of Ca\(^{2+}\) release and values of concentrations were usually saved at 0.1-ms intervals.
Fluo-3 Concentrations, Fluo-3 Fluorescence, and the Point Spread Function of the Confocal Microscope

In this article, we assume that fluo-3 can exist in four different states (Fig. 1 B): free (Fluo), bound to Ca\(^{2+}\) (CaFluo), bound to protein (PrFluo), and bound to both Ca\(^{2+}\) and protein (CaPrFluo). The fluorescence of Ca\(^{2+}\)-free fluo-3 is small whereas that of fluo-3 bound to Ca\(^{2+}\) is large and assumed to be independent of whether the complex is bound to protein. It is useful to define a new variable FFluo,

\[
[FFluo] = [CaFluo] + [CaPrFluo] + \frac{F_{min}}{F_{max}} [Fluo] + \frac{F'_{min}}{F'_{max}} [PrFluo].
\]  

(A2)

\(F_{max}\) denotes the relative fluorescence intensity of CaFluo and CaPrFluo; \(F_{min}\) and \(F'_{min}\) denote the relative fluorescence intensities of Fluo and PrFluo, respectively (Table I). According to this definition, FFluo represents the same fluorescence as the mixture of all four components.

I). According to this definition, FFluo represents the fluorescence intensities of Fluo and PrFluo, respectively (Table F).

\(\max\) denotes the relative fluorescence intensity of CaFluo and protein (CaPrFluo). The fluorescence of Ca\(^{2+}\) and protein is independent of whether the free fluo-3 is small whereas that of fluo-3 bound to Ca\(^{2+}\) is large and assumed to be independent of whether the complex is bound to protein. It is useful to define a new variable FFluo,

\[
[FFluo] = [CaFluo] + [CaPrFluo] + \frac{F_{min}}{F_{max}} [Fluo] + \frac{F'_{min}}{F'_{max}} [PrFluo].
\]  

(A2)

\(F_{max}\) denotes the relative fluorescence intensity of CaFluo and CaPrFluo; \(F_{min}\) and \(F'_{min}\) denote the relative fluorescence intensities of Fluo and PrFluo, respectively (Table I). According to this definition, FFluo represents the same fluorescence as the mixture of all four components. Consequently, fluorescence is linearly related to FFluo.

The spatial resolution of a laser scanning confocal microscope can be described by its point-spread function \(PSF\). In this article, \(PSF\) is assumed to be the product of three Gaussian functions, one each for \(x, y,\) and \(z\),

\[
PSF = \frac{1}{S_x \cdot S_y \cdot S_z} \exp \left( \frac{x^2 + y^2}{S_x^2} + \frac{z^2}{S_z^2} \right).
\]  

(A3)

The resolution factors \(S_x\) and \(S_y\) are expected to be equal since the resolution of the microscope in the plane of the image is independent of \(x\) and \(y\), consequently \(S_x = S_y = S_\omega\). The value of \(S_\omega\) is, in general, different from that of \(S_{xy}\); \(PSF\) has been normalized so that its volume integral over all space is unity.

For a Gaussian function, the distance parameter \(S_\omega\) (with \(w = x, y,\) or \(z\)) is related to the variance \(\text{var}_w\) by \(S_\omega^2 = 2\text{var}_w\). The full width at half maximum of the function, denoted by \(FWHM_\omega\), is related to \(S_\omega\) by \(FWHM_\omega = 2S_\omega (\ln 2)^{1/2}\). In the calculations described in this article, a value of 0.2 \(\mu m\) was used for \(FWHM_x\) and \(FWHM_y\), and 0.5 \(\mu m\) was used for \(FWHM_z\) (Table I).

The value of fluorescence \(F(x, y, z, t)\) produced by FFluo and detected at \(x, y, z\) is given by the convolution of \([FFluo]\) with \(PSF\),

\[
F(x, y, z, t) = \frac{A}{S_x \cdot S_y \cdot S_z} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} [FFluo](x', y', z')
\]

\[ \cdot \exp \left( -\frac{(x-x')^2 + (y-y')^2}{S_x^2} \right) \]

\[ \cdot \exp \left( -\frac{(z-z')^2}{S_z^2} \right) \]

\[ dx' dy' dz' \].

(A4)

\(A\) is a constant that depends on several factors. These include the intensity of the laser excitation light that illuminates the spot at \(x, y, z\), the absorbance of the four different states of fluo-3 and their quantum efficiencies for fluorescence, and the efficiency with which the resulting fluorescent light is collected by the microscope and measured by the photodiode(s).

Since the value of \([FFluo]\) depends on \(r' = \sqrt{x'^2 + y'^2 + z'^2}\) and not on the individual values of \(x', y',\) and \(z'\), it is natural to use spherical polar coordinates with the origin at the center of the calcium source. In this coordinate system, the values of \(x', y',\) and \(z'\) are given by

\[
x' = r' \sin \theta' \cos \phi',
\]

\[
y' = r' \sin \theta' \sin \phi',
\]

\[
z' = r' \cos \theta',
\]

and the volume element \(dx' dy' dz'\) is replaced by \(r'^2 \sin \theta' \, dr' \, d\theta' \, d\phi'\). In spherical polar coordinates, Eq. A4 takes the form

\[
F(x, y, z, t) = A \int_{0}^{\infty} W(r', x, y, z)[FFluo](r', t) 4\pi r'^2 \, dr',
\]

in which \(W(r', x, y, z)\) is given by

\[
W(r', x, y, z) = \frac{1}{4S_x S_y S_z} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{2\pi} \exp \left( -\frac{(r' \sin \theta' \cos \phi' - x)^2 + (r' \sin \theta' \sin \phi' - y)^2}{S_x^2} \right)
\]

\[ \cdot \exp \left( -\frac{(r' \cos \theta' - z)^2}{S_z^2} \right) \, d\theta' \, d\phi' \].

Thus,

\[
W(r', x, y, z) = \frac{1}{4S_x S_y S_z} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{2\pi} \exp \left( -\frac{(r'^2 \sin^2 \theta' + x^2 + y^2)}{S_x^2} \right).
\]

\[ \cdot \exp \left( -\frac{(r' \cos \theta' - z)^2}{S_z^2} \right) \, d\theta' \, d\phi' \exp \left( \frac{2r' \sin \theta' (x \cos \phi' + y \sin \phi')}{S_x} \right) \cdot \frac{2\pi r'}{S_x} \cdot \frac{2\pi r'}{S_y} \cdot \frac{2\pi r'}{S_z}. \]

(A5)

The integral with respect to \(d\phi'\) can be evaluated from equation 3.915(4) in Gradshteyn and Ryzhik (1965),

\[
\int_{0}^{2\pi} \exp \left( \frac{2r' \sin \theta' (x \cos \phi' + y \sin \phi')}{S_x} \right) \, d\phi' = 2\pi r' \frac{2r' \sin \theta' (x \cos \phi' + y \sin \phi')}{S_x^2}.
\]

(A6)
in which \( l_0 \) represents the modified Bessel function of order 0. Eqs. A5 and A6 can be combined to give

\[
W(r', x, y, z) = \frac{1}{2S_{h,y}^2 \pi^{3/2}} \int_0^\infty d\theta' \sin \theta' L \left( \frac{2\sin \theta' \sqrt{x^2 + y^2}}{S_{h,y}} \right) \exp \left( -\frac{r'^2 \sin^2 \theta' + x^2 + y^2}{S_{h,y}^2} \right) \cdot \exp \left( -\frac{(r' \cos \theta' - z')^2}{S_z^2} \right). \tag{A7}
\]

Values of \( W(r', x, y, z) \) were obtained from Eq. A7 by numerical integration.

**Normalization of \( \Delta F \) and \( \Delta [F\text{Fluo}] \) by their Resting Values**

In spark experiments, fluorescence measurements are usually expressed in terms of \( \Delta F/F_R \), in which the subscript \( R \) denotes the resting state and \( \Delta F = F - F_R \). From Eq. A4,

\[
F_R = \frac{A}{S_{h,y}^2 \pi^{3/2}} \int_0^\infty \int_0^\infty \int_0^\infty [F\text{Fluo}]_r \exp \left( -\frac{(x-x')^2 + (y-y')^2 + (z-z')^2}{S_{h,y}^2} \right) dx' dy' dz'.
\]

Because

\[
\frac{1}{S_{h,y}^2 \pi^{3/2}} \int_0^\infty \int_0^\infty \int_0^\infty \exp \left( -\frac{(x-x')^2 + (y-y')^2 + (z-z')^2}{S_{h,y}^2} \right) dx' dy' dz' = 1
\]

for finite \( x, y \), and \( z \),

\[
F_R = A[F\text{Fluo}]_r.
\]

If \( \Delta [F\text{Fluo}] \) is used to represent \( [F\text{Fluo}] - [F\text{Fluo}]_r \), the following relation, which does not depend on the value of \( A \), is obtained

\[
\frac{\Delta F}{F_R}(x, y, z, t) = \frac{1}{S_{h,y}^2 \pi^{3/2}} \int_0^\infty \int_0^\infty \int_0^\infty \frac{\Delta [F\text{Fluo}](x', y', z', t)}{[F\text{Fluo}]_r} \exp \left( -\frac{(x-x')^2 + (y-y')^2}{S_{h,y}^2} \right) \cdot \exp \left( -\frac{(z-z')^2}{S_z^2} \right) dx' dy' dz'.
\]

Since \( [F\text{Fluo}] \) depends on \( r' \) and not individually on \( x', y' \), and \( z' \), spherical polar coordinates can be introduced to give

\[
\frac{\Delta F}{F_R}(x, y, z, t) = \int_0^\infty W(r', x, y, z) \frac{\Delta [F\text{Fluo}](r', t)}{[F\text{Fluo}]_r} \cdot 4\pi r'^2 dr', \tag{A8}
\]

in which \( W(r', x, y, z) \) is given by Eq. A7.

In the spark simulations described in the first section of this Appendix and in this article, Eq. A8 is replaced by the summation

\[
\frac{\Delta F}{F_R}(x, y, z, t) = \sum_{i=0}^{100} W_i(x, y, z) \frac{\Delta [F\text{Fluo}](r(i), t)}{[F\text{Fluo}]_r}
\]

in which

\[
W_i(x, y, z) = \int_{r(i-1)}^{r(i)} W(r', x, y, z) 4\pi r'^2 dr',
\]

with \( r(-1) \) set to 0 \( \mu \)m.

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