Electric control of magnetism at room temperature

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In the single-phase multiferroics, the coupling between electric polarization (P) and magnetization (M) would enable the magnetoelectric (ME) effect, namely M induced and modulated by E, and conversely P by H. Especially, the manipulation of magnetization by an electric field at room-temperature is of great importance in technological applications, such as new information storage technology, four-state logic device, magnetoelectric sensors, low-power magnetoelectric device and so on. Furthermore, it can reduce power consumption and realize device miniaturization, which is very useful for the practical applications. In an M-type hexaferrite SrCo₂Ti₂Fe₈O₁₉, large magnetization and electric polarization were observed simultaneously at room-temperature. Moreover, large effect of electric field-controlled magnetization was observed even without magnetic bias field. These results illuminate a promising potential to apply in magnetoelectric devices at room temperature and imply plentiful physics behind them.

The magnetoelectric coupling between the magnetic and electric orders in multiferroics brings about novel physical mechanism and offers the opportunity for attractive applications, such as non-volatile memory technology, four-state logic device and magnetoelectric sensors. Of particular interest is electric control of magnetism, which has a promising potential for reducing power consumption and realizing device miniaturization in spintronic devices. However, such applications have been limited by weak magnetoelectric effects and low operating temperature in most discovered single-phase multiferroics. Very recently, giant magnetoelectric effects have been reported in some hexaferrites, in which the electric polarization (P) can be induced by their complex magnetic orders.

Hexaferrites, which have long been utilized in permanent magnets and microwave devices, can be classified into different types such as M, W, X, Y, U, and Z. Their crystal structures can be described as stacked sequences of the basic blocks: S, R, and T. Recently, some significant results about the magnetically induced ferroelectricity and resultant magnetoelectric effects were observed in some Y-, Z- and M-type hexaferrites due to the modifications of spiral magnetic structures by an external magnetic field. This phenomena can be understood by the spin current or the inverse Dzyaloshinskii-Moriya models described by the equation:

\[ P = A e_{ij} \times (\mathbf{S}_i \times \mathbf{S}_j) \]  \hspace{1cm} (1)

\((e_{ij}\) denotes spins at neighboring sites i and j, \(e_{ij}\) denotes the propagation vector of a spiral structure, \(A\) denotes the coupling coefficient determined by the exchange interaction and the spin-orbit interaction,\), in which the interaction between noncollinear spins can produce the local P through the spin-orbit interaction.\(^{(16-18)}\)

Recently, Z-type hexaferrite SrCo₂Ti₂Fe₈O₁₉ was reported to show an excellent room-temperature magnetoelectric property at low magnetic field (H), which would shed light on the practical applications of magnetoelectric multiferroics. Since the electric resistivity in most of these hexaferrites is not high enough to support an electric field, the manipulation of magnetization by an electric field has rarely been reported, especially at room temperature. Here, we report the discovery of electric control of magnetism in an M-type hexaferrite, SrCo₂Ti₂Fe₈O₁₉ (SCTFO), at room temperature. Furthermore, without magnetic bias field, large effect of electric field-controlled magnetization is observed at room temperature, which would revolutionize the field of low-power magnetoelectronic devices.

**Results**

The M-type hexaferrites, AFe₁₂O₁₉ (A = Ba, Sr, Pb, etc.), crystallize in the magnetoplumbite-type structure (the space group P6₃/mmc) with Fe³⁺ ions in octahedral (12k, 4f and 2a), tetrahedral (4f) and bipyramidal (2b) sites, respectively. This lattice structure could be symbolically described as an alternate stacking of 4 building blocks,
namely S, S*, R, and R*, which is schematically shown elsewhere15,16. The large A"+ ion, which causes a slight perturbation in the lattice around itself, is responsible for the large magnetic uniaxial anisotropy in these hexaferrites, where the crystallographic c-axis is the magnetic easy axis14. The substituted M-type hexaferrites have been extensively investigated for the application point of view15,16. Among them, the substitution of Fe"+ cations with Co"+-Ti"+ pair is reported to alter the magnetocrystalline anisotropy and give rise to the noncollinear magnetic structures20,21. In this case, Co"+ cations prefer to occupy the octahedrally-coordinated 4f and 12k Wyckoff positions and the magnetic anisotropy changes gradually from an axial anisotropy (along the hexagonal c axis) to a planar anisotropy (a, b plane of the hexagonal cell) with increasing concentration of Co2+. On the other hand, nonmagnetic Ti"+ cations substituted on the 12k would interrupt the magnetic interactions in a way that the two blocks RS and R*S* become more or less magnetically decoupled20. As a result, an intermediate conical magnetic structure was obtained in the M-type hexaferrite with critical Co-Ti concentration, which can be verified by neutron diffraction and Mössbauer spectroscopy20,23.

Figure 1 shows the room-temperature powder X-ray diffraction pattern and Rietveld refinement of polycrystalline SCTFO. All peaks can be indexed as the magnetoplumbite-type cell in the space group P63/mmc (a=5.881(1) Å, b=5.881(1) Å, c=23.102(8) Å, α=90.00°, β=90.00°, γ=120.00°). Table I summarizes the fractional coordinates and occupation factors.

The magnetic hysteresis loops for SCTFO was measured at 300 K. As shown in Fig. 2(a), the sample is ferromagnetic at room temperature with the magnetization of about 30 emu/g with \( H = 15 \) kOe. Comparing with SrFe12O19, the coercivity of SCTFO decrease remarkably (see in the inset of Fig. 2(a)), which is resulted from the reduction in anisotropy field with the substitution of Fe3+-Ti4+ for Fe3+–Ti4+. The temperature dependence of magnetization (M-T) from 100–800 K for SCTFO is shown in Fig. 2(b). The Curie temperature of 740 K for this ferrite is less than that of SrFe12O19, Fe3+-Ti4+ and Sr3Co2Fe24O41, show no spontaneous magnetization anomaly at around 420 K is likely to correspond to the transition from collinear ferrimagnetic phase to the phase with a conical magnetic structure20,22,24.

Figure 3(a) shows the temperature dependence of P at zero magnetic field, which is obtained from integrating the magnetoelectric current by time with 400 K as the initial point. As we know, the reported Y- or Z-type multiferroics hexaferrites, such as Ba3Mg2Fe12O22 or Sr3Co2Fe24O41, show no spontaneous P at zero magnetic field. When an external magnetic field is applied, a transverse conical spin structure is induced, which could produce P through the spin current and the inverse Dzyaloshinskii-Moriya mechanism11,13,18. However, in the case of SCTFO, it shows spontaneous P at zero magnetic field, indicating that its spin structure has a transverse cycloidal component, which should be further investigated by neutron diffraction measurements. In addition, a sudden decrease of P almost coincides with the anomaly in the M-T curve, suggesting the coupling between the ferromagnetism and ferroelectricity13,14. Figure 3(b) shows the H dependence of P for the sample at 300 K. With increasing H, the P is considerably suppressed, demonstrating a strong magnetoelectric coupling in SCTFO. As shown in Fig. 2(b), P does not vanish even H = 16 T, indicating that the conical spin structure is robust. Accordingly, there is no anomaly in M-H curve (Figure 3(b)), which is consistent with the behavior of P.

Table I | The fractional coordinates and occupation factors for polycrystalline SCTFO.

| Atom | Site | x    | y    | z    | SOF |
|------|------|------|------|------|-----|
| Sr1  | 2d   | 0.3333 | 0.6667 | 0.25 | 1.0 |
| Fe1  | 2a   | 0.0   | 0.0   | 0.0  | 1.0 |
| Fe2  | 2b   | 0.0   | 0.0   | 0.6231(16) | 0.5 |
| Fe3  | 4f   | 0.6667 | 0.3333 | 0.02714(28) | 1.0 |
| Co1  | 4f   | 0.6667 | 0.3333 | 0.2310(8) | 1.0 |
| Fe5  | 12k  | 0.1845(4) | 0.3691(8) | 0.11117(12) | 0.6667 |
| Ti1  | 12k  | 0.0973(5) | 0.1946(11) | 0.11266(25) | 0.3333 |
| O1   | 4e   | 0.0   | 0.0   | 0.1225(17) | 1.0 |
| O2   | 4f   | 0.3333 | 0.6667 | 0.0935(6) | 1.0 |
| O3   | 6h   | 0.8707(6) | 0.7413(13) | 0.25 | 1.0 |
| O4   | 12k  | 0.8462(9) | 0.6924(18) | 0.03472(20) | 1.0 |
| O5   | 12k  | 0.5387(6) | 1.0774(13) | 0.1144(24) | 1.0 |

Figure 1 | The room-temperature powder X-ray diffraction pattern and Rietveld refinement for polycrystalline SCTFO.

Figure 2 | (a) The magnetic hysteresis loop for SCTFO at 300 K; the inset of (a) shows the enlarged magnetic hysteresis loop; (b) The variation of M for the sample as a function of T.
Comparing with magnetic field, it is easier to generate an electric field in devices and working with a voltage, rather than an electric current, gives rise to less energy consuming in technological applications\textsuperscript{5,8,9}. Therefore, the electric-field controlled ferromagnetism, which is usually defined as converse magnetoelectric effect, is of great importance from the application point of view. However, the investigation of electric control of magnetism has been hampered in many single-phase multiferroics due to their low resistivity\textsuperscript{8,9,25}. In the case of SCTFO, the resistivity reaches up to $4.36 \times 10^{10} \, \Omega \cdot \text{cm}$, which is high enough to support an electric field and ensures the following measurement. We first measured the room-temperature magnetic loops of SCTFO with DC electric field ($E_{\text{dc}}$) parallel and perpendicular to the $H$, respectively. By applying an $E_{\text{dc}}$, an obvious reduction of $M$ is observed and typical loops at zero and 22 kv/cm field are shown in Fig. 4(a). Similar to the effect of $H$ on $P$, $M$ decreases with increasing $E_{\text{dc}}$ and the change of $M$ with $E_{\text{dc}}=22$ kv/cm reaches up to 6.3% and 2.8%, for $E_{\text{dc}} \perp H$ and $E_{\text{dc}} \parallel H$, respectively, as shown in Fig. 4(b).

The variation of $M$ of SCTFO as a function of $E_{\text{dc}}$ was measured without magnetic bias field. Before this measurement, the sample was pre-magnetized in a field of 5 kOe. A considerable reduction of $M$ is observed in Fig. 5(a) despite of positive or negative $E_{\text{dc}}$. This result is of interest since no $H$ is utilized in electric control of magnetism. Figure 5(b) shows the time dependence of $M$ under a square wave electric field of 18.5 kV/cm. $M$ decreases or increases with applying or removing $E_{\text{dc}}$, respectively, indicating a stable response to the electric field, which demonstrates the modulation of $E$ on $M$ in SCTFO.

In order to further investigate the effect of $E$ on $M$, the scanning probe microscopy (SPM) experiment was performed using lift mode. Figure 6(a) is the topography of a selected area, while Fig. 6(b)–6(d) show the magnetic force microscopy (MFM) images for the same scanning area, which reveal magnetic contrast between the surface magnetization domains. During the MFM studies, various electric fields were applied on the sample. In this case, the bias electric field

Figure 3 | (a) The temperature dependence of $P$ (red) and $M$ (black) for SCTFO. (b) The magnetic initial curve (black) and $H$ dependence of $P$ (red) at 300 K.

Figure 4 | (a) The room temperature magnetic hysteresis loop under different $E_{\text{dc}}$ (zero and 22 kv/cm); (b) The change rate of $M$ under various $E_{\text{dc}}$ with magnetic bias field ($H=460$ Oe) at room temperature: $E \perp H$ (Red); $E \parallel H$ (Black).

Figure 5 | (a) The change of $M$ as a function of $E_{\text{dc}}$ at room temperature without magnetic bias field; (b) The time dependence of $M$ under a square wave electric field.
would make the magnetic force signal shift uniformly, but would not affect the contrast of the signal. From Fig. 6, no change in domain structure can be detected, but the magnetic contrast fades out as electric voltage increasing from 0 to 10 V, indicating the effect of E on the magnetic domain. This experimental result is significant since it gives a direct and powerful evidence of electric control of magnetism at room temperature.

Discussion

In the single-phase SCTFO, the doping of Co $^{2+}$ is propitious to change the magnetic anisotropy from axial to planar, while the substitution of Ti $^{4+}$ interrupts the magnetic interactions between the two blocks RS and R$^{*}$S$^{13,20,21}$. As a result, the magnetic frustration at the boundary stabilizes the nonlinear spiral spin structure, inducing the spontaneous ferroelectric polarization at room temperature, which is attributed to the spin current or the inverse Dzyaloshinskii-Moriya models. In present sample, the substitution of a Co $^{2+}$-Ti $^{4+}$ pair can maintain the average $3 +$ valence state and then suppress the hopping of electrons between Fe $^{3+}$ and Fe $^{2+}$ in SCTFO. In addition, sintering in oxygen is also helpful to increase the resistivity. As a result, SCTFO has a high enough resistivity which ensures the magnetoelectric measurements. As shown in Fig. 3(b), the values of P decrease monotonously with the increasing magnetic field. According to Eq. (1), the spin cone in SCTFO tends to close with a magnetic field, leading to the decrease of P correspondingly$^{14}$. As for the effect of E on M, the values of M decreases with increasing E, which is obviously shown in Fig. 4, 5, and 6. When an electric field is applied on the sample, the electric polarization would intrinsically increase. According to Eq. 1, the value of $S_1 \times S_2$ would increase and lead to the enlargement of the spin cone angle correspondingly, which results in the reduction of magnetization in our sample. Since the electric control of magnetism is of significance, the effect of E on M without magnetic bias field is emphasized in present work. As shown in Fig. 5, the remanent magnetization of SCTFO can be reversibly and reproducibly controlled by an electric field, suggesting the potential application in the non-volatile rewritable RAM with multiferroic hexaferrites.

Thus we have prepared a single-phase multiferroics SCTFO. The substitution of Co $^{2+}$-Ti $^{4+}$ pair for two Fe $^{3+}$ cations produces a spiral conical magnetic structure and high resistivity in SCTFO. This special spin structure can induce spontaneous P via spin-orbit coupling, suggesting the strong magnetoelectric coupling in SCTFO. Electric control of magnetism is confirmed by magnetization and MFM measurements. The excellent magnetoelectric performance in SCTFO gives a valuable contribution to the study of magnetoelectric effect and enables the realization of electric control magnetization in magnetoelectric devices.

Methods

Poly crystalline sample of SrCo$_2$Ti$_2$Fe$_9$O$_{19}$ was prepared through conventional solid-state reaction. The high-purity powders of SrCO$_3$, Fe$_2$O$_3$, Co$_2$O$_3$, and TiO$_2$ were mixed stoichiometrically and sintered at 1,200 °C for 10 hours for twice in the oxygen atmosphere. The powder X-ray diffraction measurement was carried out to identify the sample as single phase. For the electric polarization measurements, a specimen was cut with a dimension of 3×3×1 mm$^3$ and gold electrodes were painted on both sides as a capacitance. The temperature or magnetic field dependence of P was obtained by integrating the magnetoelectric current, where temperature and applied H were controlled by the physical property measurement system (PPMS, Quantum Design) and magnetoelectric current was gathered by an electrometer (model 6514, Keithley). The magnetic properties of SrCo$_2$Ti$_2$Fe$_9$O$_{19}$ were measured with a commercial vibrating sample magnetometer (VSM, model 7407, Lakeshore). The effect of $E_0$ on $M$ was measured with a VSM under zero magnetic field at room temperature. During measurement, $E_0$ was applied on the capacitance and the induced M was detected by pick up coils. The MFM measurements were carried out on the polished sample by scanning probe microscopy (SPM) (Dimension V, Veeco). In order to investigate the effect of E on the magnetic domain, various voltages were applied between the probe tip and the sample chuck. During the MFM measuring processes, two different modes, “grounding the tip” and “applying the voltage on the tip”, are used.

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Author contribution

L.Y.W., prepared the sample; L.Y.W. and D.H.W. designed the experiments and wrote the paper. L.Y.W., H.C.X. and J.L.G. performed the measurements. L.Y.W., D.H.W., Q.Q.C., and Y.X.Z. collected and analyzed the data. All authors discussed the results and commented on the manuscript.

Additional information

Competing financial interests: The authors declare no competing financial interests.

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