Doping evolution of the electronic specific heat coefficient in slightly-doped La$_{2-x}$Sr$_x$CuO$_4$ single crystals

Seiki Komiya and I Tsukada
Central Research Institute of Electric Power Industry, Yokosuka, Kanagawa 240-0196, Japan
E-mail: komiya@criepi.denken.or.jp

Abstract. Detailed doping dependence of the electronic specific heat coefficient $\gamma$ is studied for La$_{2-x}$Sr$_x$CuO$_4$ (LSCO) single crystals in the slightly-doped regime. We find that $\gamma$ systematically increases with doping, and furthermore, even for the samples in the antiferromagnetic (AF) regime, $\gamma$ already acquires finite value and grows with $x$. This suggests that finite electronic density of states (DOS) is created in the AF regime where the transport shows strong localization at low temperatures, and this means the system is not a real insulator with a clear gap even though it still keeps long range AF order.

1. Introduction
It is still an issue of controversy in the study of high-$T_c$ cuprates how carriers are doped into the parent Mott insulator and the metallic state arises with doping. Transport measurements have demonstrated that metallic in-plane resistivity is immediately realized at high temperatures with 1%-hole doping, although it turns insulating at low temperatures[1]. This behavior is qualitatively the same even in the underdoped superconducting regime when superconductivity is suppressed by magnetic field or Zn doping[2, 3, 4]. High temperature Hall coefficient ($R_H$) measurements have shown that the temperature dependences of $R_H$ in the LSCO system can be fitted as thermal activation type at high temperatures[5], suggesting that the carrier concentration would be temperature dependent. The background electronic structure of this peculiar transport is extensively studied by angle resolved photoemission spectroscopy (ARPES) experiments, and some kind of fragmented Fermi surface, “Fermi arc”, is found to grow with doping[6, 7]. However, the identity of this novel electronic state is not fully understood yet. There is also a long debate on the metallicity of the doped Mott insulators[8, 9].

To elucidate the doping evolution of the electronic states in a different light, we study the detailed doping dependence of the electronic specific heat coefficient $\gamma$ of LSCO in slightly doped regime. We find that for the parent insulating La$_2$CuO$_4$ (LCO), $\gamma$ is indeed zero, but it systematically increases with doping even though the system keeps long range antiferromagnetic order (0.005 $\leq x < 0.02$). This result means that the finite electronic density of states is created already with 0.5%-hole/Cu doping.
2. Experimental

Single crystals of LSCO are grown by traveling solvent floating zone method\[10\]. The nominal Sr concentrations of the measured crystals are 0, 0.005, 0.01, 0.015, 0.02, 0.03, 0.04, and 0.05. These LSCO crystals in slightly-doped regime tend to have excess oxygen that leads the samples showing minor superconductivity, so these are annealed in pure Ar to remove excess oxygen. Samples with $x = 0$ to 0.015 show Néel transition and $T_N$s of these samples are determined by susceptibility measurement to be 320 K, 280 K, 240K, and 200 K for $x = 0$, 0.005, 0.01, and 0.015, respectively. Samples with typical weight of 20 mg are used for heat capacity measurement which is performed by the relaxation method, using Quantum Design’s Physical Properties Measurement System down to 2 K. Transport properties are measured by the conventional four-terminal method.

3. Results and discussions

In Fig. 1, we show the $C_p/T$ vs. $T^2$ plot of measured samples at low temperatures. We see that the low temperature specific heat gradually and systematically increases with doping, and for all samples studied here, $C_p/T$ has almost linear dependence on $T^2$ and Schottky anomaly is hardly observed in this temperature range. Thus we can reasonably fit the data in a form of $\gamma T + \beta T^3$. Note that the transition to the Mott insulating state in high-$T_c$ cuprates is not of the diverging effective mass type\[11, 12\], and the $\gamma$ values are not enhanced when the system approaches the parent Mott insulator.

The doping dependence of $\gamma$ is presented in Fig. 2. For the parent LCO, $\gamma$ is zero within an experimental error. This result is consistent with the high temperature Hall measurement. The temperature dependence of $R_H$ for the parent LCO is found to be well fitted as a thermal

![Figure 1. $C_p/T$ vs. $T^2$ plot for slightly doped LSCO crystals. Data for $0 \leq x \leq 0.02$ and $0.02 \leq x \leq 0.05$ are separately shown in the main panel and the inset, respectively.](image1)

![Figure 2. Doping dependence of the electronic specific heat coefficient.](image2)
activation type with two different activation energies[5], so γ = 0 is reasonable as an insulator with a definitive gap. When Sr is doped, γ immediately grows even in the AF “insulating” regime. This is not consistent with ARPES results which suggest a full gap in the AF regime[13]. It was also pointed out, however, that the length of the “arc” is in good correlation with the γ value in a wide doping range[7,14], and therefore similar arc with finite length which is too small to be detected by ARPES measurements would exist in the AF regime.

In Fig. 2, it is also seen that the doping dependence of γ changes at x = 0.02 where the long range AF order disappears. To see this difference in more detail, we compare the behavior of low temperature resistivity across x = 0.02. As presented in Figs. 3 and 4, ρ_{ab} for the x = 0.03 crystal shows well-known variable range hopping behavior, while for the x = 0.015 sample, low temperature resistivity can be best fitted as exp(T^{−1/2}) in a rather wide temperature range. This temperature dependence suggests that a gap, possibly a soft Coulomb gap[15], is opening with decreasing temperature for LSCO with x = 0.015. This behavior is reasonable because the effective Coulomb interaction increases when the number of carriers decreases. DOS at the Fermi level would be suppressed by this gap, causing the x dependence of γ weaker in the AF regime. Note that the finite γ value suggests that this gap may not fully open and DOS still remains. ARPES experiments would detect this gap opening.

How do the finite DOS and the AF order coexist at 0.005 ≤ x < 0.02 ? One possible scenario is related to the domain structure observed in neutron scattering experiments and the spin glass behavior at low temperatures in the AF regime[16,17]. They have found that in the AF LSCO samples with 0.01 ≤ x < 0.02, some part of the AF ordered phase turns into a spin glass phase at low temperatures. The observed γ in the AF regime would come from this spin glass portion in the background AF order. Magnetotransport experiments support this domain-structure scenario, where large negative magnetoresistance suddenly appears above a magnetic field of the weak ferromagnetic transition[18].

Although the transport properties show the strong localization in slightly to underdoped regime at low temperatures[1,2], we emphasize that the system is not an insulator with a gap in the energy spectrum. This can be also understood by comparing the behavior of the
thermal conductivity and the heat capacity. In LSCO, low temperature thermal conductivity measurements have revealed that the residual electronic term of the thermal conductivity at zero temperature limit is zero in the non-superconducting samples with $0 \leq x \leq 0.05^{[19]}$, while the electronic specific heat coefficient increases monotonically with $x$ in the same doping range. Therefore, the observed resistivity divergence is simply because the mean free length of carriers decreases with decreasing temperature. Although the localization mechanism itself is still unclear, the electronic ground state may not be different regardless of whether the long range AF order is established.

4. Summary
Low temperature electronic specific heat is studied for the LSCO system in and near the AF regime. $C_p$ follows the simple $\gamma T + \beta T^3$ law at low temperatures, and electronic specific heat coefficient $\gamma$ is found to increase systematically with doping even in the AF regime ($0 \leq x < 0.02$). This result suggests that a finite density of states at the Fermi level is already created with 0.5%-hole doping, and that the doped Mott insulator is not a real insulator with a clear gap.

References
[1] Ando Y, Lavrov A N, Komiya S, Segawa K and Sun X F 2001 Phys. Rev. Lett. 87 017001
[2] Boebinger G S, Ando Y, Passner A, Kimura T, Okuya M, Shimoyama J, Kishio K, Tamasaku K, Ichikawa N and Uchida S 1996 Phys. Rev. Lett. 77 5417 – 20
[3] Ono S, Ando Y, Murayama T, Balakirev F F, Betts J B and Boebinger G S 2000 Phys. Rev. Lett. 85 638 – 41
[4] Komiya S and Ando Y 2004 Phys. Rev. B 70 060503(R)
[5] Ono S, Komiya S and Ando Y 2007 Phys. Rev. B 75 024515
[6] Yoshida T et al 2003 Phys. Rev. Lett. 91 027001
[7] Yoshida T et al 2007 J. Phys.: Condens. Matter 19 125209
[8] Lee P A, Nagaosa N and Wen X G 2006 Rev. Mod. Phys. 78 17 – 85
[9] Choy T P and Phillips P 2005 Phys. Rev. Lett. 95 196405
[10] Komiya S, Ando Y, Sun X F and Lavrov A N 2002 Phys. Rev. B 65 214535
[11] Kumagai K, Suzuki T, Taguchi Y, Okada Y, Fujishima Y and Tokura Y 1993 Phys. Rev. B 48 7636 – 42
[12] Padilla W J, Lee Y S, Dunn M, Blumberg G, Ono S, Segawa K, Komiya S, Ando Y and Basov D N 2005 Phys. Rev. B 72 060511(R)
[13] Shen K M et al 2004 Phys. Rev. B 69 054503
[14] Momono N, Ido M, Nakano T, Oda M, Okajima Y and Yamaya K 1994 Physica C 233 395 – 401
[15] Efros A L and Shklovskii B I 1975 J. Phys. C 8 L49 – 51
[16] Matsuda M, Fujita M, Yamada K, Birgeneau R J, Endoh Y and Shirane G Phys. Rev. B 65 134515
[17] Niedermayer Ch, Bernhard C, Blasius T, Golnik A, Moodenbaugh A and Budnick J I 1998 Phys. Rev. Lett. 80 3843 – 46
[18] Ando Y, Lavrov A N and Komiya S 2003 Phys. Rev. Lett. 90 247003
[19] Takeya J, Ando Y, Komiya S and Sun X F 2002 Phys. Rev. Lett. 88 077001