A STUDY OF DYNAMICAL PROCESSES IN THE ORION KL REGION USING ALMA—PROBING MOLECULAR OUTFLOW AND INFLOW

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ABSTRACT

This work reports high spatial resolution observations toward the Orion KL region with high critical density lines of CH3CN (124–114) and CH3OH (8–1, 8–70, 7), as well as a continuum at ∼1.3 mm band. The observations were made using the Atacama Large Millimeter/Submillimeter Array with a spatial resolution of ∼1.5′ and sensitivity of about 0.07 K and ∼0.18 K for continuum and line, respectively. The observational results showed that the gas in the Orion KL region consists of jet-propelled cores at the ridge and dense cores east and south of the region that are shaped like a wedge ring. The outflow has multiple lobes, which may originate from an explosive ejection, and is not driven by young stellar objects. Four infrared bubbles were found in the Spitzer/IRAC emissions. These bubbles, the distributions of the previously found H2 jets, the young stellar objects, and molecular gas suggest that BN is the explosive center. The burst time was estimated to be ≤1300 yr. At the same time, signatures of gravitational collapse toward Source I and the hot core were detected with material infall velocities of 1.5 km s⁻¹ and ∼0.6 km s⁻¹, corresponding to mass accretion rates of 1.2 × 10⁻³ M⊙/yr and 8.0 × 10⁻⁵ M⊙/yr, respectively. These observations may support the belief that high-mass stars form via the accretion model, similar to their low-mass counterparts.

Key words: ISM: jets and outflows – ISM: kinematics and dynamics – stars: individual (Orion BN/KL, Source I) – stars: pre-main sequence

Online-only material: color figures

1. INTRODUCTION

The dynamical processes in the massive star formation region are still not well understood. Molecular outflow and inflow are critical for testing whether or not high mass stars form in the same way as their low mass counterparts. The Orion BN/KL is the massive star formation region closest to us. The Orion KL region consists of different spatial components including the “ridge” and “core.” The velocity ranges of these components are from 2.5 to 9 km s⁻¹ (Wright et al. 1996 and the references therein). Every component harbors high-mass young stellar objects (YSOs) identified with near- and mid-infrared emissions. Many of the YSOs have subcomponents revealed by high spatial resolution observations (Shuping et al. 2004 and the references therein).

The first high-velocity molecular outflow was identified and discovered using the CO J = 1–0 line in the Orion KL region (Kwan & Scoville 1976). The outflow attracted extraordinary attention and was further imaged by dense molecular tracers, such as HCO*J = 1–0, and SO (3–2–1) (Schulz et al. 1995; Friberg 1984, and the references therein). Besides the high-velocity phenomena detected in millimeter line emissions, finger-like H2 filaments were also discovered in the BN/KL region by near-infrared observations (Taylor et al. 1984). The line emissions of atoms, ions, and high energy level CO were detected in the outflow region too, indicating the presence of shocked and hot gas in the outflow (Doi et al. 2002; Gry et al. 1999; Beuther et al. 2010). Different masers including H2O, OH, SiO, and CH3OH were identified or imaged in the Orion KL outflow region, and the masers may be excited by shocks (Matthews et al. 2010; Hirota et al. 2012 and the references therein). The properties of the molecular outflow are still under investigation. Recently, Zapata et al. (2012) observed an outflow in the Orion KL region using the SiO (J = 8–7, v = 0) line with the Atacama Large Millimeter/Submillimeter Array (ALMA). The high-velocity gas lobes are shaped like a butterfly. The driving source of the outflow in the KL region remains uncertain. It may be produced by the disruption of the massive stellar system, which occurred ∼500 yr ago and ejected massive stellar objects BN, I, and n (Zapata et al. 2009; Bally et al. 2011). While Beuther & Nissen (2008) suggested that the submillimeter core SMA1 may host the driving source of the outflow. Source I was identified as the driving source of the SiO (J = 8–7, v = 0) outflow (Zapata et al. 2012). The origin of the outflow needs to be further tested and explored.

Material infall toward cores or forming stars is another kind of essential motion in star formation regions. Although the infall studies progress more slowly than those of molecular outflows, gravitational collapse candidates were frequently found toward massive star formation regions with a blue profile (Wu & Evans 2003; Wu et al. 2007; Chen et al. 2010, and the references therein). The stronger collapse signature, inverse P Cygni profile was detected early with high resolution observations (Welch et al. 1987; Zhang & Ho 1997; Qin et al. 2008; Liu et al. 2013, and the references therein). The dynamical collapse and inflow motion in this active star formation region were not reported before.

To examine the dynamical processes and especially to uncover the origin of the outflow in this region, observations with tracers of high excitation temperature and critical density are needed. This paper presents the results of an observational study of Orion KL at 1.3 mm with ALMA. Using the continuum...
emission and spectral lines of CH$_3$CN (12$_2$–11$_1$) and CH$_3$OH (8–70) in both emission and absorption, an outflow with multiple lobes in the KL region was detected. Inflow motion was identified toward Source I and the hot core.

2. OBSERVATIONS

The data used in this paper were obtained from ALMA science verification (SV) at band 6. The observations were made with 16 × 12 m antennas of ALMA on 2012 January 20 with baselines ranging from 18 to 253 kλ. The phase center was at R.A. = 05h35m14.35 and decl. = −05°22′35′′. There are 20 spectral windows (SPW) of 1.875 GHz wide. Each spectral window consists of 3840 channels with a spectral resolution of 0.488 MHz (∼0.7 km s$^{-1}$). Callisto was used for bandpass calibration and flux calibration, while quasar J0607-085 provided the phase calibration. The primary beam size of FWHM is about 30″ and the emission of the KL region falls well inside the FWHM.

The synthesized images with calibrated data were carried out using the Common Astronomy Software Applications (CASA) package and MIRIAD (Sault et al. 1995). The continuum map was constructed from the line-free channels of the five spectral windows at highest frequency bands. The natural-weighted beam size and the rms noise level for the continuum are 1′′66 × 1′′35 with P.A. = −20°7 and 8 mJy beam$^{-1}$ (0.07 K), respectively. The CH$_3$CN (12$_2$–11$_1$) and CH$_3$OH (8–70) lines were extracted from spectral windows “spw1” and “spw19,” respectively, while the $^{13}$CH$_3$CN (12$_2$–11$_1$) and (13$_3$–12$_2$) transitions were identified in spectral windows “spw5” and “spw6,” respectively. The typical rms for lines is about 20 mJy beam$^{-1}$ per channel (0.18 K).

3. RESULTS

3.1. Continuum

The continuum image of the 1.3 mm emission is presented in Figure 1(a). Seven millimeter cores are identified and denoted with MM followed by a number. MM1 is the strongest and largest, and corresponds to the hot core (Wilson et al. 2000). The others are located in the compact ridge. MM7 is too far away from the phase center and thus excluded from further analysis. The known infrared sources were plotted on the map and are shown by crosses (Shuping et al. 2004). Two-dimensional Gaussian fits were made toward the continuum cores. The peak position is listed in Column 2 of Table 1. Column 3 is the core size; here 414 pc was adopted as the distance (Menten et al. 2007). The peak emission intensity and the total flux of the cores are given in Columns 4 and 5 of Table 1.

$^{13}$CH$_3$CN (12$_2$–11$_1$) and (13$_3$–12$_2$) transitions were extracted from each of the cores. The systematic velocities ($V_{lsr}$), and the mean line widths ($\Delta V$) of the $^{13}$CH$_3$CN lines by Gaussian fitting are given in Columns 6 and 7 of Table 1. We also calculated the rotational temperatures of the cores with the rotational temperature diagram (RTD) method (Araya et al. 2005; Goldsmith & Langer 1999; Liu et al. 2002; Qin et al. 2010). Only the lines not blended and with high S/N levels (>5σ) are used for RTD analysis. Assuming that the local thermodynamic equilibrium (LTE) holds and the lines are optically thin, $T_{rot}$ values are obtained and are listed in the eighth column of Table 1. Figure 1(b) shows the rotational temperature diagram of MM6 as an example. We estimated the virial mass of the MM cores with the relation (MacLaren et al. 1988): $M_{vir} = 210(R/pc)(AV/km s^{-1})^2(M_\odot)$ which is given in the last column of Table 1.

![Figure 1](image-url)
3.2. Line Profiles and Spatial Distribution of gas Emissions

Figure 2 shows the spectral lines of CH$_3$CN (12--11) and CH$_3$OH (8--7) from the peak position of each dust core observed with ALMA. There are two kinds of line profiles from these cores. One has the emission distributed in both sides of the systemic velocity (e.g., the lines at MM1 and MM2 positions). Both CH$_3$CN and CH$_3$OH lines at MM1 have double peaks, with the blue peak stronger than the red one, and the absorption dip is red shifted with respect to the systemic velocity. These absorption dips are unlikely to be caused by missing flux, as the instrument effect cannot account for the absorption, which can be interpreted as gas inflow motion toward the core (see more in Section 3.4).

The other kind of line profiles shows that most of the gas emission is located in one side of the systemic velocity. The spectral line at MM4 has stronger emission on the blue side, and on the red side the intensity decreases sharply from the peak. In contrast to the MM4, the spectral emission of MM3, MM5, and MM6 is mainly on the red side. These profiles show characteristics of high-velocity gas moving toward us or away from us, which are similar to the fast moving molecular bullets near YSOs as shown in Figures 7 and 9 of Bachiller et al. (1990), but there are apparent differences. The high-velocity gases observed here are not molecular bullets ejected from the immediate vicinity of the YSOs, but are jet propelled cores (hereafter JPCs). First, the JPCs in Orion KL show strong dust emission with total flux ranging from 0.43 to 1.27 Jy, which is very rare among the bullets in bipolar molecular outflows or HH objects. Their virialization masses range between those of molecular bullets in bipolar outflows near the YSOs.

Second, the molecular bullets near YSOs are always highly collimated, and the blue and red bullets in high- or low-mass star formation regions are located along an outflow axis (Richer et al. 1992; Hatchell et al. 1999). In the Orion KL region the JPCs have no such bipolar structure, but appear to distribute within the inner edge of the gas region. Third, the JPCs are hot with $T_{\text{rot}}$ of 100 to 182 K. Line profiles of bullets in bipolar molecular outflows have usually been observed in the low transition lines of CO $J = 1-0$, 2--1 (Bachiller et al. 1990; Hatchell et al. 1999). However, for the JPCs in Orion KL, no such profile was found in the lines of CO $J = 1-0$, 3--2, HCO$^+$ $J = 1-0$, and SO (3--2) (Ericksen et al. 1982; Olofsson et al. 1982; Friberg 1984). In addition, the spectral profiles of the JPCs are different from those of molecular bullets in bipolar outflows near the YSOs. The former have a single component, whereas the latter are usually accompanied by a line from the host molecular core (Bachiller et al. 1990; Richer et al. 1992; Hatchell et al. 1999), indicating that the JPCs in the Orion KL region were not ejected by a YSO. Their systemic velocity is not separated from the high-velocity emission, meaning that the original gas in the KL region was impacted and propelled by the explosive jets (Taylor et al. 1984; Bally et al. 2011; Zapata et al. 2011). These cores are not explosive jets, but jet-propelled cores.

The spectral features can also be seen from channel maps of the CH$_3$CN (12--11) line illuminated in Figure 3, where the white contours denote the emission intensity averaging over 3 km s$^{-1}$ velocity range. The gas emission at velocity $\lesssim$(-13) km s$^{-1}$ is contaminated by the CH$_3$CH$_2$CN (25--24--24) line. The major bulk of the gas is distributed within the velocity range of $-5$ to 16 km s$^{-1}$. A comparison with Figure 1(a) shows that as the velocity increases, the emission first arises from the west of MM1, and then from MM4 and southeast of MM5. From 1 to 16 km s$^{-1}$, the morphology of the emission region looks like a semi-wedge ring. A complete wedge ring appears with a velocity ranging from 6 to 11 km s$^{-1}$, which is similar to an angle open to the north. These results show that the molecular gas emissions of the KL...
region are physically related to different velocity components, including molecular JPCs and cores.

In the blue-shifted, high-velocity range of $-14$ to $-2$ km s$^{-1}$, gas emission arises from MM4, the southeast of MM5, and the west of MM1, and then extends to the northeast and south. The red-shifted, high-velocity gas emissions arise from MM6 to MM5, the southwest of MM1, and near MM2/MM3, with the velocity range of $-19$ to $22$ km s$^{-1}$. The components of the high-velocity gas also consist of the emissions of JPCs and cores.

### 3.3. High-velocity Gas—an Outflow with Multiple Lobes

From Figure 2, the blue- and red-shifted, high-velocity gases range from $-15$ to $-3$ km s$^{-1}$ and $15$ to $24$ km s$^{-1}$, respectively. The contours of the velocity integration intensities of the high-velocity gas measured with the CH$_3$CN (12$_1$–11$_1$) line are presented in Figure 4, in which the crosses represent the infrared sources.

The CH$_3$CN (12$_1$–11$_1$) outflow has two blue lobes and five red lobes. The blue-shifted, high-velocity gas peaks near MM4 and to the west of MM1. The red-shifted, high-velocity gas reaches maximum near MM5. The spectra of MM1 and Source I have a core emission profile with blue and red high-velocity wings, but the lines of the JPCs have much more obvious blue- or red-shifted emissions (Figures 2 and 5; see also Section 3.4). The distribution of the high-velocity gas suggests that the outflow mainly overlaps or connects with JPCs. The dissimilarities between the spectra of the cores and the JPCs also demonstrate that the molecular outflow is not from MM1 or Source I.

The red outflow has five lobes and looks like a projection of a bracelet, whereas the blue outflow has two lobes. The outflow covers JPCs MM4–MM6. The southwest part of MM1 overlaps the outflow, and Source I is located outside the outflow. Neither MM1 nor Source I can be the origin or the driving source of the outflow. The other YSOs in this region are located in different lobes and cannot serve as the driving source of the outflow either (see Figure 4). These results demonstrate that the high-velocity gas in the KL region is not caused by a usual bipolar molecular outflow derived by a YSO, but is mainly from distinct molecular cores propelled by bullets from a burst. That is why blue and red lobes cannot be separated, and the bipolar outflow detected in the KL region with CO at the early period has poor collimation (Kwan & Scoville 1976; Wu et al. 2004).

### 3.4. Gravitational Collapse toward Source I and Inflow Motion to the MM1

To examine the line emissions of MM1 (the hot core) and Source I, we plot the CH$_3$OH (8–1, $v$=7) and $^{13}$CH$_3$CN (13$_3$–12$_3$ lines at the boundaries of MM1 and Source I in Figure 5. The CH$_3$OH spectral profile from Source I shows blue-shifted emission and red-shifted absorption with respect to the systematic velocity of 5.5 km s$^{-1}$. It presents an inverse P Cygni profile, indicating that both the gases behind and in front of the continuum source are moving toward the continuum source.

The spectrum of CH$_3$OH at MM1 has a gas self-absorption dip that is red-shifted relative to the systematic velocity. The blue peak of the CH$_3$OH line is stronger than the red one with a ratio of $T_B/T_R = 2.16$. For the line profile, the asymmetry parameter $\delta V (= V_{\text{thick}} - V_{\text{thin}})/\delta V_{\text{thin}}$ defined by Mardones et al. (1997) was calculated, giving $\delta V \leq -0.25$ for the blue profile, where $V_{\text{thick}}$ is the peak velocity of the optical thick line, and $V_{\text{thin}}$ and $\delta V_{\text{thin}}$ are the systemic velocity and the line width of the optically thin line, respectively. For the line pairs of the center of MM1, $\delta V = (V_{\text{thick}} - V_{\text{thin}})/\delta V_{\text{thin}} = -0.26$ was obtained. The two parameters of $T_B/T_R$ and $\delta$ both satisfy the criteria of the blue profile.

The optical thick and thin spectral lines from position (d), east of MM1, have similar intensity and velocity partner to those from the center of MM1. The ratio of $T_B/T_R = 2.13$ and the $\delta V = -0.25$ satisfy the criteria of a blue profile. The signature is slightly weaker than that of the center one, suggesting that inflow motion in MM1 is moving toward the core center.

The line of the position (e), from MM1, has a profile that belongs to those of JPCs and has blue-shifted, high-velocity
gas. This suggests that the west side of the core MM1 was impacted by the explosive jets. The profiles of the lines from north and south of MM1 seem to be affected by the explosive outflow, but not as significantly as the one at position (e).

4. DISCUSSION

4.1. The JPCs

As stated before, the molecular gas of the KL region actually consists of different velocity components. Beside the cores, the JPCs are mainly distributed in the compact ridge. There are two kinds of line profiles produced from the JPCs: one has blue wings and the other has red wings. The high-velocity wings in the two spectra are from $-13$ to $-3$ km s$^{-1}$ and 15 to 24 km s$^{-1}$, respectively, which are similar to those from SiO (8–7, $v = 0$) (Zapata et al. 2012). Such high velocities may be excited by C-type shock because they are below the critical velocities of J-type shocks, which range from 25 to 45 km s$^{-1}$. J-type shocks are also found with emissions primarily in the UV and optical, whereas C-type shocks mainly excite the vibrational and high rotational levels of molecules in infrared and radio emissions (Wang & Scoville 1992). Unlike the spectra of molecular bullets near the YSOs, the emissions of ambient and high-velocity gas of the JPCs are not separated. From Table 1 and Figure 5, one can see that the $V_{lsr}$ of the JPCs are different. The spectra of JPCs with blue wings have smaller $V_{lsr}$, whereas those with red wings have larger $V_{lsr}$. If we take the average value of the peak velocities of all the JPCs and cores in the KL region, 7.7 km s$^{-1}$ as the systemic velocity of the KL region gas, then the shifts of the peak velocities of the JPCs range from $-2.4$ to 2.3 km s$^{-1}$. The JPCs are moving away or toward us at a slow speed. The slow, systemic motions of the JPCs may be caused by a rocket effect when the JPCs were propelled by the jets.

4.2. Explosive Molecular Outflow

4.2.1. High-velocity Molecular Outflow Driven by Explosive Jets

The high-velocity molecular outflow has multiple lobes (see Figure 4). The outflow gas mainly distributes in the ridge and only a small part overlaps with the southeast of MM1 and the north of MM2, indicating that the high-velocity gas comes mainly from the JPCs. From Figures 2 and 5, one can see that the blue-shifted JPCs are located east of the ridge and the red ones distribute more widely. Unlike bipolar molecular outflows in which a YSO is usually driving the high-velocity gas emission, every blue or red lobe in the KL region harbors YSOs. Sources I and IRn are located far from the outflow center. These results suggest that the outflow is unlikely to be driven by some young stellar objects in the outflow region.

To examine the origin of the high-velocity gas and shock effects on the JPCs, we investigated the distribution of the IRAC [4.5/3.6] flux ratio in the Orion KL region. The 4.5 $\mu$m band emissions detected with IRAC of Spitzer contain lines excited by shocks (Cyganowski et al. 2009). A ratio $>1.5$ represents shocks or jets (Takami et al. 2010). The ratio map, with a color scale overlaid with the continuum map, is displayed in Figure 6. Four infrared bubbles (IRBs) were identified, and the largest and brightest one, IRB1, is centered at BN and the emissions are saturated at the inner region. IRB2 is the second largest and is connected with IRB1 and IRB3. IRB4 is located south of the region. Among the ALMA-detected dust cores, MM6 and MM5 are located at the outer part of IRB1 and IRB2, respectively. MM4 is located at the outer part of IRB 3. The northeast of MM3 is connected with IRB4. These bubbles have the flux ratio between $\lesssim 1.0$ and 1.4, and seem to be swept out by wind and jets. The JPCs are located outside the bubbles, or partially overlap with bubbles. It is notable that the strongest and largest shocked area in the KL region is a large wedge ring.
and surrounds the KL gas wedge. This ring is well associated with the H$_2$ 2.12 and 2.15 μm clumps (Stolovy et al. 1998), suggesting shocked effects in this region.

To further probe the excitation mechanism of the molecular outflow, we also plot the JPCs and cores of the KL region, overlaid with the 2.12 μm H$_2$ jets (Bally et al. 2011 and the references therein) in Figure 7. Clearly the JPCs do not overlap with the finger-like jets. The western side of the core MM1, where the shocked gas (Figure 5) was detected, partially overlaps with a finger-like jet. These results indicate that the jets break into the region and the cores are propelled or impacted. The gas is accelerated and entrained, which results in the multi-lobe molecular outflow.

We identify the BN object as the explosive center in the next section. Here we estimate the dynamical age of the expanding bubble. Using the distance from BN to the project center of the high-velocity gas as the average length of the jets (8′′), and using the average velocity of the outflow, the dynamic time is estimated to be 1280 yr, which is consistent with the numerical experiments of gradational slingshot (Chatterjee & Tan 2012).

4.2.2. The Explosive Center

One astonishing characteristic seen from Figures 6 and 7 is the distribution of YSOs. Almost all the YSOs in the ridge and BN region are located at the edges or outside the jets, whereas all the YSOs are located within the bubbles or on the rings. BN is located at the center of the largest bubble IRB1. IRc15, IRc16W, IRc16S, IRc16N, IRc17, IRc6E, IRc6, and IRc6S are located in IRB1 (also see Shuping et al. 2004, Figure 1). IRc3N and IRc3S are embedded in the second largest bubble IRB2. IRc2, IRc7, and IRc8 are located in IRB3. IRc4 is located in IRB4. Recent evidence of triggered star formation was found in a number of bubbles (Watson et al. 2008; Deharveng et al. 2003). The ionizing stars responsible for the bubbles are located in the wind-evacuated cavity. A neutral material layer forms between the ionization front and the shock front when the H$^+$ region expands. This layer may collapse, fragment, and then form stars (Elmegreen & Lada 1977; Churchwell et al. 2007; Watson et al. 2008; Deharveng et al. 2003). The shocked layers can be produced by various mechanisms, including colliding clouds, expanding H$^+$ regions, stellar wind bubbles, and supernova remnants. For the Orion BN/KL region, there are explosive jets, whereas the BN has a hypercompact H$^+$ region (Plambeck et al. 2013). The facts about the gas distribution and the position of BN in the bubble suggest that BN is very likely the explosive or expanding center.

BN is the hottest and evolved one among the stellar sources in the region. It is a B-type star with emission variation (Becklin & Neugebauer 1967). At the wavelengths 12.5 μm and shorter infrared wavelengths, BN is the dominating object (Shuping et al. 2004 and the references therein).

In addition to the BN, there are many YSOs in the hot core and ridge. These YSOs were identified from near or middle infrared observations around the BN at different times and with different sensitivities. Their numbering generally has no order. Although the first five YSOs discovered were numbered with IR1-IR5 located southeast of BN (IR1) (Rieke et al. 1973), the sources found later cannot be numbered following this way. For example, IR6 and IR7 were detected and surrounded by the BN, IR2, and IR3 (Downes et al. 1981). Because of the large extinction at near- and mid-infrared wavelengths, their positions may be not so accurate. However, the following analysis showed that the evolutionary states of these YSOs are correlated with their distance from the BN object.

The YSOs of IRc6, IRc3, and IRc 20 are closest to BN, and are likely more evolved than those in other places of the region. Their mid-infrared emission are rather weak (Shuping et al. 2004) and strong emissions at longer wavelengths were not seen. At the northwest fringe of the wedge ring, there is no cold gas emission detected with the inversion lines of $^{15}$NH$_3$ (3,3) (Wilson et al. 2000), meaning that IRc6 is really hot.

The objects IRc2, IRn, and IRc7 are located farther away from BN and were younger than those of IRc6, IRc3, and IRc20. IRc2 A+B, IRc7, and IRcn have much stronger mid-infrared emissions than those of IRc6, IRc3, and IRc20 in the IRB2. VLA observations showed that a shell of water masers was...
located at IRc2 A, which is the connected region of the infrared cluster IRc2 and MM1 (Plambeck & Wright 1988).

IRc4 is the farthest from the BN among the YSOs in the ridge. Infrared spectroscopic observations show that IRc4 is separated from IRc2 (Aitken et al. 1981). It has little emission at middle infrared, but has strong emission at 20 μm wavelength (Shuping et al. 2004 and the references therein). More recently, SOFIA observations show that IRc4 has the strongest emission at 37.1 μm (De Buizer et al. 2012), suggesting that it is the youngest among the YSOs in this region.

The physical state of the gas in the Orion-KL region shows that the BN is the power source. Figure 3 shows that the cavity of the hot gas wedge ring detected with CH3CN (12–11) is open toward the BN object, as inferred from the observed direction. From Figure 6 one can see that the largest shock ring clearly opens to the BN object. MM6, MM5, and MM3 are located along the axis of the wedge. Such a structure can also be seen in images of other wavelengths. Infrared continuum measurements at 2–30 μm with a resolution 2″–2′2 show that the KL region is like a clumpy cavity with a diameter of about 10″ (Wynn-Williams et al. 1984). It is also shaped like a wedge ring. The most northern infrared emission is located in its west fringe and covers IRc6. However in this region there is no emission of the rather cold gas detected with the inversion line of 15NH3 (3,3), meaning the gas closest to the BN is hot.

The evolutionary states of the YSOs in the ridge correlate with their distances from the BN object, which also suggests that the BN is the powering source of the KL region.

Various mechanisms were suggest for the explosive center of the Orion BN/KL region (see Section 1). In this work we report some observational characteristics for BN as the explosive center. More theoretical analysis and observation are needed to confirm its identification.

4.3. Dynamics of Source I and the Hot Core MM1

The inverse P Cygni profile of line (b) shown in Figure 5 indicates gravitational collapse toward Source I. The material infall velocity can be estimated as 1.5 km s⁻¹ from the V_red and V absorption of the absorption peak (Welch et al. 1987). Here we take the V_red of Source I as same as that of the core MM1. Assuming the infall is free fall (Pineda et al. 2014), the mass accretion rate of 1.2 × 10⁻³ M⊙/yr can be obtained with dM/dt = 3V²_red/2G, which is much larger than those of low mass cores (Myers et al. 1996). The result is consistent with the mass of Source I, which is about 10 M⊙ (Testi et al. 2010; Liu et al. 2013).

For the core MM1, the lines at the center and position (d) have blue profiles showing inflow motion. The infall velocity is calculated with the model of Myers et al. (1996):  

\[ V_{in} \approx \sigma^2 \left( \frac{\sigma_{blue}}{\sigma_{red} - \sigma_{blue}} \right) \frac{1 + e^{T_{BD}/T_D}}{1 + e^{T_{RD}/T_D}} \]  

where \( T_D \) is the brightness temperature of the dip (assumed to be optically thick), \( T_{BD} \) and \( T_{RD} \) are the height of the blue and red peaks above the dip, respectively. The velocity dispersion \( \sigma \approx 2.7 \) km s⁻¹ is obtained from the optically thin \(^{13}\)CH3CN(13–12) line. Thus the infall velocities inferred from CH3CN (12–11) and CH3OH (8–70) are \( \sim 0.4 \) and \( \sim 0.8 \) km s⁻¹, respectively. Taking the average value \( 0.6 \) km s⁻¹ as \( V_{in} \), the mass accretion rate of 8.0 × 10⁻⁵ M⊙/yr was obtained. Both the infall velocity and the mass accretion rate are smaller than those of Source I, but still larger than those of low mass cores (Myers et al. 1996).

Signatures of gravitational collapse or inflow motion were detected in the Orion KL region for the first time (see Section 1). For Source I, the material infall velocity and mass accretion rate are computable with those of high mass sources (see Liu et al. 2013 and the references therein). Early, an outflow from the Source I was detected with \( v = 0, J = 2–1 \) line of SiO (Plambeck et al. 2009), which is from northeast to southwest and is mainly distributed within the hot core. Around Source I, an ionized disk was imaged with radio and 7 mm continuum emission (Reid et al. 2007). Recently a hot and neutral circumstellar disk around Source I was detected with ALMA (Hirota et al. 2014). The collapse signature, disk, and outflow suggest that the high-mass YSO in Source I forms similarly to those of low-mass stars.

Blue profiles were detected in the hot core, indicating that the inflow motion is weaker than that in Source I. Previous observations showed that inflows occur more frequently and more strongly in UCHII regions than in UCHII precursors (see Wu et al. 2007 and the reference therein). The inflow difference of Source I and MM1 is attributed to the fact that Source I is a radio source, while no YSO has been found in the hot core yet.

5. CONCLUSION

We have studied Orion KL region with 1.3 mm continuum and lines of CH3CN (12–11), \(^{13}\)CH3CN(13–12), and CH3OH (8–70) observed with ALMA. Seven dust cores, MM1 to MM7, were detected. Emissions of IRAC 3.6 and 4.5 μm were included in the analysis. Our main findings are as follows:

1. Molecular emissions of CH3CN (12–11) and CH3OH (8–70) were detected from the dust cores. The derived rotation temperatures \( T_{rot} \) are from 119 to 198 K. Lines from MM1 and MM2 have the usual profiles of cores, and those from MM3, MM4, MM5, and MM6 show jet propelled line characteristics. The line from the western area of MM1 also presents the JPC profile similar to that of MM4. The gas emission of the Orion KL region consists of these JPCs and the cores. The morphology of the major body of the gas emission is similar to a wedge ring.

2. The high-velocity gas ranges from \( -13 \) to \( -3 \) km s⁻¹ and \( 15 \) to \( 24 \) km s⁻¹ for the blue and red wings, respectively. The morphology of the outflow shows multiple lobes. The high-velocity gas mainly comes from the JPCs and overlaps with MM1 at its southwest and with the north of MM2. The outflow does not seem to be driven by YSOs like usual bipolar outflows, but instead originates from an explosion.

3. Four infrared bubbles were found with the ratio of [4.5]/[3.6] from infrared emissions with IRAC and Spitzer, which are related to molecular bullets and contain the YSOs in the region. The properties of the BN object, the possible formation mechanism, and evolutionary states of the YSOs, as well as the distribution of molecular gas in this region, were analyzed. Results seem to support that the BN object is the explosive center and the explosion occurred 1300 yr ago.

4. The properties and spectral profile of Source I and core MM1 are presented. An inverse P Cygni profile was found toward Source I, showing gravitational collapse. The infall velocity and mass infall rate are 1.5 km s⁻¹ and 1.2 × 10⁻³ M⊙/yr, respectively. Inflow motion shown in the blue profile was detected toward MM1 (hot core), which has
infall velocity 0.6 km s\(^{-1}\) and a mass infall rate \(8.0 \times 10^{-5} \, M_\odot/\text{Yr}\).

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