Induction experiments have been performed for nearly one hundred years for measuring the hardness of materials.\(^1\) Recent years have seen increased interest in indentation because of the significant improvement in indentation equipment and the need for measuring the mechanical properties of materials on small scales. It is now possible to monitor, with high precision and accuracy, both the load and displacement of materials during indentation. It can, in principle, depend on parameters, which is sometimes difficult, especially when "piling-up" occurs.

Using a scaling approach to indentation problems developed recently,\(^7,8\) we reveal, in this letter, an approximate relationship between the ratio of hardness to elastic modulus and the ratio of irreversible work to total work in indentation. Together with a well-known relationship between elastic modulus, initial unloading slope, and contact area, a new method is then suggested for estimating the hardness and modulus of solids using instrumented indentation with conical or pyramidal indenters. \(\odot 1998\) American Institute of Physics. [S0003-6951(98)02931-3]

\(\sigma = K \varepsilon^n,\) for \(\varepsilon > Y/E,\)

where \(K\) is the strength coefficient and \(n\) is the work-hardening exponent.\(^9\) To ensure continuity, \(K = Y(E/Y)^n\). Consequently, either \(E, Y,\) and \(K\) or \(E, Y,\) and \(n\) are sufficient to describe the stress-strain relationship. We use the latter set of parameters extensively in the following discussions. When \(n\) is zero, Eq. (1) becomes the model for elastic-perfectly plastic solids. For most metals, \(n\) has a value between 0.10 and 0.50.\(^{10}\)

During loading, the force, \(F,\) on the indenter can be written, according to dimensional analysis, as\(^7,8\)

\[F = E h^2 \Pi_a (Y/E, \nu, n, \theta),\]

where \(\Pi_a\) is a dimensionless function of four dimensionless parameters \(Y/E, \nu, n,\) and the indenter half angle \(\theta.\) From Eq. (2), the total work done by the indenter, \(W_{\text{tot}},\) to cause elastic and plastic deformation when the indenter reaches a maximum depth, \(h_m,\) is given by

\[W_{\text{tot}} = \int_0^{h_m} F dh = \frac{E h_m^3}{3} \Pi_a \left(\frac{Y}{E}, \nu, n, \theta\right).\]

Thus, the force and work are proportional to \(h_m^2\) and \(h_m^3,\) respectively.

Since unloading takes place after loading during which the indenter reaches the maximum depth, \(h_m,\) the force, \(F,\) on the indenter can be written, according to dimensional analysis, as\(^7,8\)

\[F = E h^2 \Pi_b \left(\frac{Y}{E}, \frac{h}{h_m}, \nu, n, \theta\right),\]

where \(\Pi_b\) is a dimensionless function of five dimensionless parameters, \(Y/E, h/h_m, \nu, n,\) and \(\theta.\) The force is, in general, no longer proportional to the square of the indenter displacement. It can, in principle, depend on \(h/h_m\) through the di-
mensionless function $\Pi_k$. Furthermore, the final depth, $h_f$, at which the force on the indenter first becomes zero during unloading, can be obtained by solving

$$F = Eh_m^3 \Pi_k \left( \frac{Y}{E} \frac{h_f}{h_m}, v, n, \theta \right) = 0.$$  \hspace{1cm} (4')

Clearly, $h_f/h_m$ is a function of $Y/E$, $v$, $n$, and $\theta$.

The work done by the solid to the indenter during unloading, $W_u$, can now be expressed as

$$W_u = \int_{h_f}^{h_m} Fdh = Eh_m^3 \int_1^1 x^2 \Pi_k \left( \frac{Y}{E} x, v, n, \theta \right) dx = Eh_m^3 \Pi_k \left( \frac{Y}{E} x, v, n, \theta \right).$$  \hspace{1cm} (5)

where the integral, $\int_{h_f/h_m}^{h_m} x^2 \Pi_k (Y/E, x, v, n, \theta) dx$, is obviously a dimensionless function, $\Pi_k$, of $Y/E$, $v$, $n$, and $\theta$. Consequently, the work done on the indenter during unloading is also proportional to $h_m$. The ratio of irreversible work to total work for a complete loading-unloading cycle, $1 - W_u/W_{tot}$, can be written as

$$\frac{W_{tot} - W_u}{W_{tot}} = 1 - \frac{\Pi_k (Y/E, v, n, \theta)}{\Pi_k (Y/E, v, n, \theta)} = \Pi_k \left( \frac{Y}{E} x, v, n, \theta \right).$$  \hspace{1cm} (6)

This ratio is, therefore, independent of $h_m$.

Because there is, as yet, no analytical solution to the problem of conical indentation in elastic-plastic solids, we chose to use finite element calculations to evaluate the dimensionless function, $\Pi_k$. Since there are only a few indenter angles of experimental interest, we fix $\theta = 68^\circ$ as in early studies (e.g., Ref. 11) and evaluate the dependence on $Y/E$, $v$, and $n$ to reveal the essential physics of the work of indentation. The details of the finite element model have been given previously. In Fig. 1 the function $\Pi_k$ is evaluated for many values of $Y/E$, $v$, and $n$. Consistent with common experiences, the ratio of irreversible work to total work decreases with increasing $Y/E$ and $n$, i.e., when the solid is more "elastic."

In a previous study, we have examined the functional dependence of hardness $H$, defined as the force divided by the projected contact area under full load, on $E$, $v$, $Y$, and $n$. Using Eq. (18) in Ref. 8, the ratio of hardness to Young’s modulus, $E$, can be written as

$$\frac{H}{E} = \Pi_k \left( \frac{Y}{E}, v, n, \theta \right).$$  \hspace{1cm} (7)

where $\Pi_k$ is a dimensionless function of $Y/E$, $v$, $n$, and $\theta$. The value of $\Pi_k$ calculated using the same finite element method is shown in Fig. 2. As expected, the value of $H/E$ increases with increasing $Y/E$ and $n$.

While much of the physics of hardness and irreversible work of indentation are contained in Eqs. (6) and (7), it is instructive to investigate the relationship between $H/E^*$ and $1 - W_u/W_{tot}$, where $E^* = E/(1 - v^2)$. In Fig. 3, we replot all the data points shown in Figs. 1 and 2 in the form of $H/E^*$ vs $1 - W_u/W_{tot}$. It is evident from Fig. 3 that all the data points lie approximately on a single curve. Consequently, we may write

$$\frac{H}{E^*} = \Pi_k \left( \frac{W_{tot} - W_u}{W_{tot}} \right).$$  \hspace{1cm} (8)

FIG. 1. Scaling relationship for $(W_{tot} - W_u)/W_{tot} = \Pi_k (Y/E, v, n, \theta)$ for $\theta = 68^\circ$.

FIG. 2. Scaling relationship for $H/E = \Pi_k (Y/E, v, n, \theta)$ for $\theta = 68^\circ$.

FIG. 3. Relationship between $H/E^*$ and $(W_{tot} - W_u)/W_{tot}$, including data from finite element calculations for conical indents and experimental results for a few materials using pyramidal indenters. (a) For $0.0 < (W_{tot} - W_u)/W_{tot} < 1.0$ and (b) details for $0.9 < (W_{tot} - W_u)/W_{tot} < 1.0$.}

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The functional dependence of $\Pi_\theta$ is displayed in Fig. 3. The subscript $\theta$ denotes a possible dependence on indenter angle, since Eq. (8) and Fig. 3 were obtained for a particular indenter angle. Nevertheless, Eq. (8) shows that, for a given indenter angle, there is an approximate one-to-one correspondence between $H/E^*$ and $1 - W_s/W_{tot}$. Consequently, the value $H/E^*$ may be obtained from the measurement of $W_s$ and $W_{tot}$, which can be determined readily by simple numerical integration based on force and displacement measurements.

The ratio of hardness to elastic modulus, $H/E^*$, is of significant interest in both tribology and fracture mechanics. This ratio multiplied by a geometric factor is the “plasticity index” which describes the deformation properties of a rough surface in contact with a smooth surface. When the plasticity index is much less than unity, the deformation of asperities is likely to be entirely elastic. When it is significantly greater than one, the deformation is predominantly plastic. In fracture mechanics, the quantity of $H/E^*$ appears in various fracture toughness equations. A few methods have been proposed to obtain this ratio, including that by Marshall et al. from the measurement of the diagonals of Knoop indentation. The present method provides an alternative method for measuring $H/E^*$ on micro- and nanometer scale for both metals and ceramics.

In the following, we show that it is possible to obtain both $H$ and $E^*$ using Eq. (8) and the well-known relationship,

$$\frac{dF}{dh}\bigg|_{h=h_m} = \frac{2}{\sqrt{\pi}} E^* \sqrt{A}, \quad (9)$$

where $dF/dh|_{h=h_m}$ is the initial unloading slope and $A$ is the projected contact area evaluated at $h_m$. Equation (9) was initially derived for indentation into elastic solids. Using the infinitesimal theory of continuum mechanics, we have recently generalized it to the initial unloading in elastic-plastic solids using indenters with axisymmetric smooth profiles. Furthermore, we have shown analytically that Eq. (9) holds true even for materials with work hardening and initial stress.

Using Eq. (9) and the definition of hardness under full load, we obtain

$$\frac{H}{E^*^2} = \frac{4}{\pi} \frac{F}{(dF/dh)^2}, \quad (10)$$

in which $F$ and $dF/dh$ are both evaluated at $h_m$. Therefore, the ratio of $H/E^*^2$ can be obtained from instrumented indentation experiments. Consequently, both $H$ and $E^*$ can, in principle, be obtained by combining Eqs. (8) and (10), i.e.,

$$H = \frac{\pi}{4} \Pi_\theta \left( \frac{W_{tot} - W_s}{W_{tot}} \right) \left( \frac{dF}{dh} \right)^2 F^{-1}, \quad (11)$$

$$E^* = \frac{\pi}{4} \Pi_\theta \left( \frac{W_{tot} - W_s}{W_{tot}} \right) \left( \frac{dF}{dh} \right)^2 F^{-1}.$$

Several experimental data points obtained using Berkovich diamond indenters from literature for Cu, W, Al, fused silica, and sapphire are also included in Fig. 3. The finite elastic constants of the diamond indenter are taken into account by using the reduced modulus commonly defined as:

$$1/E^* = (1 - \nu^2)/E + (1 - \nu^2)/E_s,$$

where $E_s = 1140$ GPa and $\nu_s = 0.07$ are the Young’s modulus and Poisson’s ratio of the diamond indenter, respectively. Since these experimental points lie approximately on the same curve as the finite element results, Eq. (8) seems thus applicable to indentation using these pyramidal indenters. Furthermore, since Eq. (9) has been shown to hold true for the pyramidal indenters with small (<3%) corrections, it appears that Eq. (11) would, therefore, be applicable to indentation using pyramidal indenters.

We would like to re-emphasize that the relationship between $H/E^*$ and the work of indentation shown in Figs. 3(a) and 3(b) as well as Eq. (8) is an approximate one based on observation of a large number of finite element calculations and a few experimental data points. Nevertheless, it appears to be sufficient for estimating $H/E^*$ to within about 15% using the mean value at each $(W_{tot} - W_s)/W_{tot}$ shown in Fig. 3. The accuracy may be improved by noticing the work-hardening dependency especially when $H/E^*$ is small [Fig. 3(b)], provided that the work-hardening exponent, $n$, is known. Alternatively, the Oliver and Pharr procedure may be used to estimate the contact area to determine $H$ and $E^*$, since this method is reasonable for materials that moderately work harden. However, we encourage additional experiments, as well as modeling efforts, to truly ascertain the accuracy of the method.

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