BAT AGN Spectroscopic Survey (BASS) – VI. The $\Gamma_x$–$L/L_{\text{Edd}}$ relation

Benny Trakhtenbrot,1† Claudio Ricci,2,3 Michael J. Koss,1,4‡ Kevin Schawinski,1 Richard Mushotzky,5 Yoshihiro Ueda,6 Sylvain Veilleux,5 Isabella Lamperti,1,7 Kyuseok Oh,1 Ezequiel Treister,2 Daniel Stern,8 Fiona Harrison,9 Mislav Baloković9 and Neil Gehrels10§

1Institute for Astronomy, Department of Physics, ETH Zurich, Wolfgang-Pauli-Strasse 27, CH-8093 Zurich, Switzerland
2Instituto de Astrofísica, Facultad de Física, Pontificia Universidad Católica de Chile, Casilla 306, Santiago 22, Chile
3Kavli Institute for Astronomy and Astrophysics, Peking University, Beijing 100871, China
4Eureka Scientific Inc., 2452 Delmer St. Suite 100, Oakland, CA 94602, USA
5Department of Astronomy and Joint Space-Science Institute, University of Maryland, College Park, MD 20742, USA
6Department of Astronomy, Kyoto University, Kyoto 606-8502, Japan
7Astrophysics Group, California Institute of Technology, 4800 Oak Grove Drive, MS 169-224, Pasadena, CA 91109, USA
8Cahill Center for Astronomy and Astrophysics, California Institute of Technology, Pasadena, CA 91125, USA
9NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA

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ABSTRACT
We study the relation between accretion rate (in terms of $L/L_{\text{Edd}}$) and shape of the hard X-ray spectral energy distribution (namely the photon index $\Gamma_x$) for a large sample of 228 hard X-ray-selected, low-redshift active galactic nuclei (AGNs), drawn from the Swift/BAT AGN Spectroscopic Survey (BASS). This includes 30 AGNs for which black hole mass (and therefore $L/L_{\text{Edd}}$) is measured directly through masers, spatially resolved gas or stellar dynamics, or reverberation mapping. The high-quality and broad energy coverage of the data provided through BASS allow us to examine several alternative determinations of both $\Gamma_x$ and $L/L_{\text{Edd}}$. For the BASS sample as a whole, we find a statistically significant, albeit very weak correlation between $\Gamma_x$ and $L/L_{\text{Edd}}$. The best-fitting relations we find, $\Gamma_x \simeq 0.15 \log L/L_{\text{Edd}} + \text{const.}$, are considerably shallower than those reported in previous studies. Moreover, we find no corresponding correlations among the subsets of AGN with different $M_{\text{BH}}$ determination methodology. In particular, we find no robust evidence for a correlation when considering only those AGN with direct or single-epoch $M_{\text{BH}}$ estimates. This latter finding is in contrast to several previous studies which focused on $z > 0.5$ broad-line AGN. We discuss this tension and conclude that it can be partially accounted for if one adopts a simplified, power-law X-ray spectral model, combined with $L/L_{\text{Edd}}$ estimates that are based on the continuum emission and on single-epoch broad-line spectroscopy in the optical regime. We finally highlight the limitations on using $\Gamma_x$ as a probe of supermassive black hole evolution in deep extragalactic X-ray surveys.

Key words: black hole physics – galaxies: active – quasars: general – X-rays: galaxies.

1 INTRODUCTION
One of the major goals in the study of active galactic nuclei (AGN) is to understand how basic physical properties of the accreting supermassive black hole (SMBH) are linked to the emergent (continuum) radiation field. The ultraviolet (UV)-optical continuum can be explained by (thin) accretion discs, in a way which involves the BH mass ($M_{\text{BH}}$), accretion rate (in terms of the Eddington ratio, $L/L_{\text{Edd}}$), and the BH spin, through deterministic, analytical and/or numerical models (e.g. Davis & Hubeny 2006; Done et al. 2012; Netzer 2013, and references therein). This is not the case with the X-ray continuum emission, which is thought to originate from a compact, hot corona that surrounds the inner parts of the accretion disc and Compton upscatters the disc UV photons. Indeed, it is not yet clear...
whether this significant emission component can be directly linked to any key AGN properties, from both theoretical and observational perspectives.

In the energy range $\sim$0.5–10 keV, the intrinsic X-ray continuum emission is observed to follow a power law of the form $dN/dE \propto E^{-\Gamma_x}$. Early evidence for a correlation between $\Gamma_x$ and $L/E_{\text{edd}}$ was put forward by several studies that focused mainly on specific high-$\Gamma_x$ AGN and/or on narrow-line Seyfert 1 sources (e.g. Pounds, Done & Osborne 1995; Brandt, Mathur & Elvis 1997; Brandt & Boller 1998; Porquet et al. 2004; Wang, Watarai & Mineshige 2004; Bian 2005). These sources, which are generally thought to represent the high-$L/E_{\text{edd}}$ end of the (local) AGN population, exhibit soft X-ray spectra (i.e. high $\Gamma_x$). However, the limited size and range of luminosities probed in these early studies prohibited them from ruling out a scenario where the fundamental underlying relation is driven by $L_{\text{AGN}}$ (or $M_{\text{BH}}$), rather than $L/E_{\text{edd}}$.

Since then, several studies have provided an increasingly more complete picture of this proposed relation by probing AGN that cover a wide range of luminosities and redshifts. These include the studies of Shemmer et al. (2006) and Shemmer et al. (2008, hereafter S08), which were the first to provide measurements for a substantial sample of $z > 3$, extremely luminous quasars; Riallit, Young & Elvis (2009, hereafter R09), which relied on a large sample of quasars drawn from the Sloan Digital Sky Survey (SDSS), at $0 \lesssim z \lesssim 4.5$; Brightman et al. (2013, hereafter B13), which used dozens of unobscured, moderate-luminosity AGN at $0.5 \lesssim z \lesssim 2$, from the COSMOS survey; and Fanali et al. (2013), which studied a sample of unobscured AGN from the XMM-Newton Bright Serendipitous Survey. The common result of these studies is the identification of a robust, statistically significant positive correlation between $\Gamma_x$, usually measured over the observed-frame range of $\sim$2–10 keV, and $L/E_{\text{edd}}$. Moreover, these studies demonstrated that $L/E_{\text{edd}}$ is indeed the main driver of this relation, and not $L_{\text{AGN}}$ and/or $M_{\text{BH}}$.

Most recently, the study of Brightman et al. (2016) showed that the $\Gamma_x/L/E_{\text{edd}}$ relation is also applicable to local Compton-thick (CT) AGN, using a small sample of sources for which precise, maser-based determinations of $M_{\text{BH}}$ are available. Finally, several X-ray variability studies identified a trend of increasing $\Gamma_x$ with increasing flux levels for individual systems (e.g. Magdziarz et al. 1998; Pounds, Done & Osborne 1995; Brandt, Mathur & Elvis 1997; Brandt & Boller 1998; Porquet et al. 2004; Wang, Watarai & Mineshige 2004; Bian 2005). These sources, which are generally thought to represent the high-$L/E_{\text{edd}}$ end of the (local) AGN population, exhibit soft X-ray spectra (i.e. high $\Gamma_x$). However, the limited size and range of luminosities probed in these early studies prohibited them from ruling out a scenario where the fundamental underlying relation is driven by $L_{\text{AGN}}$ (or $M_{\text{BH}}$), rather than $L/E_{\text{edd}}$.

A promising way for addressing some of these limitations, and for expanding the $\Gamma_x/L/E_{\text{edd}}$ relation towards more complete, larger samples of AGN, is to study hard X-ray-selected AGN, for which the spectral coverage in the X-rays extends to higher energies. Indeed, several studies tried to identify relations between $\Gamma_x$ and $L/E_{\text{edd}}$ in samples of AGN detected by the hard X-ray Swift/BAT instrument (covering roughly 15–150 keV; Gehrels et al. 2004), which are essentially free of any obscuration-related selection biases. Some of these earlier Swift/BAT studies demonstrated the significant scatter in the $\Gamma_x–L/E_{\text{edd}}$ plane, and found no convincing evidence for a correlation between these properties – interpreted as a result of the limited sample size (e.g. Winter et al. 2009a,b). The increasing size of Swift/BAT-detected AGN samples, combined with more elaborate X-ray spectral analyses, eventually allowed the identification of significant $\Gamma_x/L/E_{\text{edd}}$ correlations (Winter et al. 2012; Kawamuro et al. 2016a; note that in the former study the correlations are found only when binning the sample by $L/E_{\text{edd}}$, similarly to B13).

In this study, we seek to establish a relation between $\Gamma_x$ and $L/E_{\text{edd}}$ for a large and essentially complete sample of low-redshift, hard X-ray-selected AGN. Our sample is based on the first data release of the BAT AGN Spectroscopic Survey (BASS). BASS provides a rich collection of X-ray and optical data for about 642 AGN, mostly at $z < 0.5$, with unprecedented levels of completeness in terms of optical spectroscopy. Compared to other low-redshift AGN samples, the hard X-ray selection that forms the basis of BASS ensures that the resulting sample is minimally affected by the AGN hosts, particularly by obscuring dust and/or contaminating optical line emission. The BASS sample covers a wide range in $L_{\text{AGN}}$, $M_{\text{BH}}$, and $L/E_{\text{edd}}$, for AGN of essentially all emission line and/or obscuration-based classification. It therefore serves as an ideal benchmark for addressing many open questions concerning the X-ray and optical emission mechanisms in AGN, and how these are related to basic BH properties (Berney et al. 2015; Lamperti et al. 2017; Oh et al. 2017).
This paper is organized as follows. In Section 2, we present our sample and the data from which we measure $L/L_{edd}$ and $\Gamma$ s. In Section 3, we examine possible correlations between these quantities, but conclude that robust and/or strong correlations of this sort cannot be clearly established for our BASS sample of AGN. In Section 4, we discuss our main findings, in the context of the several previous studies that reported $\Gamma_s-L/L_{edd}$ relations. Section 5 summarizes our findings. Throughout this work we assume a cosmological model with $\Omega_m = 0.7$, $\Omega_{\Lambda} = 0.3$, and $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$.

2 SAMPLE AND DATA

2.1 The BASS dr1 sample of AGN

This work focuses on AGN selected through their hard-band X-ray emission, as identified in the Swift/BAT 70-month catalogue (Bauermann et al. 2013). Out of 1210 unique objects in that catalogue, 836 have been identified as known AGN. The first data release of the BASS project (Koss et al., in preparation; hereafter K17) includes 642 of these AGN, for which redshifts and complementary multiwavelength data are available. As part of the BASS effort, we have curated optical spectra for 580 AGN, which were then used to measure accurate redshifts and other spectral properties, relying on narrow emission lines (usually [O iii] 5007; see K17). For 62 additional AGN, redshifts are available from the NASA Extragalactic Database (NED).

Of the initial sample of 642 sources with redshifts, we first focus on the 425 AGN for which determinations of $M_{BH}$ are available within BASS. More information regarding these determinations of $M_{BH}$ is provided in Section 2.3 below. We further focus on sources within the redshift range 0.01 < $z$ < 0.5, thus omitting 40 AGN and leaving in 385 BASS sources. This is done to avoid high-$z$-beamed AGN (and blazars), and also extremely nearby AGN for which the precise distances (and therefore luminosities) may be somewhat uncertain. We finally select only those sources for which the available X-ray observations have a sufficiently high number of counts, i.e. $N_{counts} > 1000$, to ensure a high-quality spectral fit (more information regarding our X-ray data is given in Section 2.2 below).

Our final primary sample therefore consists of 228 AGN at 0.01 < $z$ < 0.5, which have a high-quality X-ray spectrum and a reliable BH mass determination. These include 30 AGN with `direct' mass measurements (either from masers, gas or stellar dynamics, or reverberation mapping); 149 AGN with $M_{BH}$ estimates obtained through single-epoch spectra of broad Balmer lines; and 49 AGN for which $M_{BH}$ is determined by combining stellar velocity dispersion ($\sigma_v$) measurements and the $M_{BH}$-$\sigma_v$ relation. We note that – unlike other samples that investigated the relation between the BH mass and the X-ray luminosity ($L_x$) – our final sample consists of both broad- and narrow-line AGN (174 and 54 AGN, respectively).

2.2 X-ray data and analysis

The analysis of the available X-ray data for the BASS AGN was presented in detail in Ricci et al. (submitted; hereafter R17). This analysis included all the X-ray data available for the BASS sample, including Swift/XRT, XMM-Newton/EPIC, Chandra/ACIS, Suzaku/XIS, or ASCA/GIS/SIS observations, and typically covering the observed-frame energy range of 0.3–150 keV. The X-ray data were fitted with a set of models that rely on an absorbed power-law X-ray SED with a high-energy cut-off, and a reflection component. A cross-calibration constant was applied to each source, in order to account for possible flux variability between the 70-month integrated Swift/BAT spectrum and the significantly shorter 0.3–10 keV observations. Additional components accounting for warm absorbers, soft excess, Fe Kα lines, and/or other spectral features were added if deemed necessary to obtain a satisfactory fit to the data. The reader is referred to R17 for a detailed discussion of the models’ physical components, parameters, and fitting quality. We note here that the R17 analysis did not explicitly impose a finite range of possible $\Gamma_s$.

The analysis of the X-ray data provided several ways of determining $\Gamma_s$, which we use throughout this study:

(i) First, $\Gamma_{tot}$ denotes the photon index recovered from the entire (relevant) energy range and the full multicomponent model adopted for each source (in which $E_C$ is a free parameter). This is the fiducial photon index adopted in the R17 study and throughout the present work (unless otherwise noted).

(ii) $\Gamma_{0.3-10}$ results from modelling the entire X-ray spectral range with a model which ignores the high-energy cut-off (i.e. setting $E_C = 500$ keV).

(iii) $\Gamma_{0.1-1}$ denotes the photon index that describes only the observed-frame 0.3–10 keV energy range, using a model that ignores the high-energy cut-off (which was fixed to $E_C = 500$ keV) and the reflection component.

(iv) $\Gamma_{BAT}$ results from fitting a power-law model solely to the energy range probed by Swift/BAT (i.e. 14–195 keV).

Fig. 1 presents some of the statistical properties of $\Gamma_{tot}$ for our sources. The left-hand and centre panels show $\Gamma_{tot}$ and the related uncertainty ($\Delta\Gamma_{tot}$), plotted against the number of counts across the available X-ray spectral range, $N_{counts}$. As noted above, in this work we include only BASS AGN with $N_{counts} > 1000$, where the typical uncertainty on $\Gamma_{tot}$ is $\Delta\Gamma_{tot} < 0.2$. This choice, which we can make only thanks to the high-quality X-ray data in BASS, can be considered conservative – indeed, previous studies of the $\Gamma_{tot}$ relation relied on X-ray spectra with significantly fewer counts. The left-hand panel of Fig. 1 suggests that our cut on $N_{counts}$ does not bias our sample against any particular range in $\Gamma_{tot}$.

Our primary sample of 228 AGN with $N_{counts} > 1000$ and 0.01 < $z$ < 0.5 covers a wide range in $\Gamma_{tot}$ with $\Gamma_{tot} \sim 1$–2.8, a median (and mean) value of $\langle \Gamma_{tot} \rangle = 1.8$, and a standard deviation of $\sigma(\Gamma_{tot}) = 0.27$.

The R17 analysis of the X-ray data for our sources also provides a more detailed and complete method for quantifying the obscuration towards the BASS AGN, based on the (Hydrogen) column density $N_{HI}$. Setting the threshold at log ($N_{HI}$/cm$^{-2}$) = 22 splits our primary sample to 162 unobscured and 66 obscured AGN (i.e. with log [$N_{HI}$/cm$^{-2}$] below and above 22, respectively, and still obeying the redshift and $N_{counts}$ cuts described above). 27 of the AGN in our primary sample are heavily obscured, with log ($N_{HI}$/cm$^{-2}$) $\geq$ 23.5, and 8 of these are CT (log ($N_{HI}$/cm$^{-2}$) $\geq$ 24; see also Ricci et al. 2015).

2.3 Black hole masses, bolometric luminosities and $L/L_{edd}$

The BH masses available for all our 228 BASS sources were determined through several different methods. First, for 30 sources, we relied on directly measured $M_{BH}$ – either from masers; spatially resolved gas or stellar dynamics; or from reverberation mapping.

1 For example, only about $\sim$1/3 of those studied by B13 had $N_{counts} \gtrsim 1000$ (a cut of $N_{counts} > 250$ was applied for the entire B13 sample).
For 149 AGN, our $M_{\text{BH}}$ estimates rely on single-epoch spectra of the broad Balmer emission lines, and prescriptions that fundamentally rely on the results of reverberation mapping campaigns. In particular, for 126 sources, we used the broad H$\beta$ emission line and the adjacent continuum luminosity ($L_{5100} \equiv \lambda L_{\lambda}[5100 \, \text{Å}]$), relying on the same line fitting procedure and $M_{\text{BH}}$ estimator as in Trakhtenbrot & Netzer (2012). For 23 additional sources, $M_{\text{BH}}$ is determined from the broad H$\alpha$ emission line, following the procedure described in Oh et al. (2015) and the prescription of Greene & Ho (2005). Finally, for 49 sources with no broad emission lines, we used $\sigma_*$ measurements and the $M_{\text{BH}}-\sigma_*$ relation of Kormendy & Ho (2013).

As explained in the BASS/DR1 paper (K17), we prefer to use the ‘direct’ $M_{\text{BH}}$ determinations, whenever available. Otherwise, we use the single-epoch estimates from broad Balmer lines, and finally those from $\sigma_*$. This reflects the different levels of uncertainty related to each of the mass estimation methods, which are discussed in K17. We briefly note here that the uncertainties on BH masses derived through single-epoch spectra of broad lines – which constitute the largest subset in our BASS sample – may reach ~0.3–0.4 dex (see e.g. Shen & Liu 2012; Shen 2013; Peterson 2014; Mejía-Restrepo et al. 2016, and references therein). On the other hand, for $M_{\text{BH}}$ determinations based on resolved stellar or gas dynamics (including masers), the statistical uncertainties are much lower, $\lesssim$0.1 dex. Importantly, the single-epoch mass estimators are calibrated in a way that minimizes any systematic offsets with respect to other methods (see e.g. Park et al. 2012; Grier et al. 2013; Woo et al. 2013).

We estimated the bolometric luminosities of our sources, $L_{\text{bol}}$, following several different prescriptions, based on the available X-ray and optical luminosities of our AGN. We mainly use the (absorption-corrected) luminosities in the 2–10 keV rest-frame energy range, $L_{2-10}$, derived from the best-fitting, multicomponent spectral models of the X-ray data (but ignoring any cross-calibration scaling factors; see R17). These are combined with three different bolometric corrections. First, we used a fixed bolometric correction of $f_{\text{bol,2–10 keV}} \equiv L_{\text{bol}}/L_{2-10} = 20$, a typical value for AGN (see e.g. Elvis et al. 1994; Marconi et al. 2004; Vasudevan & Fabian 2007; Jin et al. 2012). Second, we used the $L_{2-10}$-dependent bolometric corrections of Marconi et al. (2004). For the sample considered in this study, these are in the range of $f_{\text{bol,2–10 keV}}$ = 11–140, with a median value of $f_{\text{bol,2–10 keV}}$ = 26.7, and 80 per cent of the sources having $f_{\text{bol,2–10 keV}} \simeq 18–48$. The resulting $L_{\text{bol}}$ are therefore slightly larger than those obtained through $f_{\text{bol,2–10 keV}} = 20$, by 0.1 dex (median value; the standard deviation is 0.18 dex), but otherwise there are no significant systematic differences between the two. We have also examined the effects that an $L_{\text{bol}}$-based bolometric correction would have on our results. For this, we relied on the results of Vasudevan & Fabian (2007), which provide $f_{\text{bol,2–10 keV}} = 20$ for $L/L_{\text{Edd}} \leq 0.04$, $f_{\text{bol,2–10 keV}} = 70$ for $L/L_{\text{Edd}} \geq 0.4$, and follow $f_{\text{bol,2–10 keV}} \propto L/L_{\text{Edd}}$ over the range $0.04 < L/L_{\text{Edd}} < 0.4$. The more recent study of Jin et al. (2012) suggests a similar dependence of $f_{\text{bol}}$. We stress that these prescriptions for $L_{\text{bol}}$ may provide markedly different values for individual sources, and therefore potentially affect any analysis of the $\Gamma_x-L_{\text{Edd}}$ plane. We indeed consider them all in our analysis (see Section 3.2).

We additionally used the absorption-corrected BAT luminosities, which cover the range 14–150 keV, combined with a fixed bolometric correction of $f_{\text{bol,14–150 keV}} \equiv L_{\text{bol}}/L_{\text{BAT}} = 8.5$. This bolometric correction is derived from the $f_{\text{bol,2–10 keV}} = 20$ one, by assuming a constant $\Gamma_x = 1.8$ – similar to the median value our sample (see Fig. 1), which corresponds to $L_{\text{BAT}}/L_{2-10} = 2.35$. These $L_{\text{BAT}}$-based estimates of $L_{\text{bol}}$ are generally in very good agreement with the fiducial, $L_{2-10}$-based ones. The median difference is 0.04 dex (with $L_{\text{BAT}}$-based estimates of $L_{\text{bol}}$ being slightly lower), and the standard deviation is 0.21 dex.

Finally, for the subset of 126 AGN for which $M_{\text{BH}}$ was determined from single-epoch spectroscopy of the broad H$\beta$ line, we derived an additional set of $L_{\text{bol}}$ estimates using $L_{\text{bol}}$-dependent bolometric corrections, $f_{\text{bol,5100}} (5100 \, \text{Å})$, which are calibrated against the Marconi et al. (2004) ones (see also Trakhtenbrot & Netzer 2012). In the range of $L_{5100}$ covered by our BASS sample, these can be approximated by...
referred to as ‘marginally significant’, in order to avoid a situation where

\[ M \sim L \]

terminations of \( \Gamma_1 \) above, and regardless of the method used for

The analysis presented in some previous studies of the

3.1 Straightforward analysis with \( \Gamma_{\text{tot}}-L/L_{\text{Edd}} \)

Fig. 2 shows the photon index versus the accretion rate for the entire (parent) sample of 425 BASS AGN. We stress that this includes all the non-blazar sources for which the quantities are available, ignoring (for now) the different redshift and \( N_{\text{counts}} \) cuts described above, and regardless of the method used for \( M_{\text{BH}} \) estimation. Here, we use the photon index we obtained from the entire spectral fit to the available X-ray data, \( \Gamma_{\text{tot}} \), and the \( L/L_{\text{Edd}} \) estimates that are based on \( L_{\text{bol}} = 20 \times L_{2-10} \).

A formal (Spearman) hypothesis test results in a weak and only marginal statistically significant correlation between the quantities, with the probability of finding a correlation if the null hypothesis (i.e. no correlation) is true being \( P = 0.8 \) per cent, and a correlation coefficient of \( r_s = 0.23 \). Thus, it appears that our parent BASS sample may hold limited evidence for a \( \Gamma_1-L/L_{\text{Edd}} \) relation of the kind found in several previous studies, although at lower statistical significance (<3\( \sigma \)). However, in what follows we will demonstrate that this result is not robust, and in particular that it does not hold for subsets of sources that differ in the \( M_{\text{BH}} \) determination methodology, for alternative determinations of \( L/L_{\text{Edd}} \), and/or when some data quality cuts are imposed on the sample.

In Fig. 3, we again show \( \Gamma_{\text{tot}} \) versus \( L/L_{\text{Edd}} \), but only for the 228 BASS AGN in our main sample, i.e. those that satisfy \( N_{\text{counts}} \geq 1000 \) and \( 0.01 < z < 0.5 \). Here, too, we use \( \Gamma_{\text{tot}} \) and the \( L_{2-10} \)-based estimates of \( L/L_{\text{Edd}} \). Fig. 3 also shows the best-fitting relations between \( \Gamma_1 \) and \( L/L_{\text{Edd}} \) reported in the three main reference studies of S08, R09, and B13. Adopting a notation of

\[ \Gamma_1 = \alpha \log (L/L_{\text{Edd}}) + \beta, \tag{1} \]

these studies have reported \((\alpha, \beta) = (0.31, 2.11), (0.31, 2.28),\) and \((0.32, 2.27)\), respectively.4 The samples and methods used in these studies are described in Section 4.2.

As Fig. 3 clearly shows, there is a considerable amount of scatter and little evidence for strong trends between \( \Gamma_{\text{tot}} \) and \( L/L_{\text{Edd}} \) in our sample of 228 BASS AGN. In an attempt to illustrate the overall trends that may be present in our sample, in Fig. 4 we show the binned \( \Gamma_1 \) versus \( L/L_{\text{Edd}} \) for each of the \( \mathcal{M}_{\text{BH}} \) subsets, where the bins spread 0.5 dex in \( L/L_{\text{Edd}} \). The markers represent the median values within each bin, while the vertical error bars represent the median absolute deviations (MAD) of \( \Gamma_{\text{tot}} \). Fig. 4 further demonstrates the large scatter in the (underlying) BASS sample, and the limited evidence for a strong \( \Gamma_1-L/L_{\text{Edd}} \) correlation for our AGN. A formal correlation test does indeed show evidence for a weak, but statistically significant correlation: the null hypothesis of no correlation between \( \Gamma_{\text{tot}} \) and \( L/L_{\text{Edd}} \) can be rejected at a level corresponding to \( P = 1.65 \times 10^{-2} \) per cent, when the entire sample of 228 AGN is considered. The corresponding Spearman correlation coefficient is \( r_s = 0.31 \) – implying a weak correlation.5 The results of this and other correlation tests are given in Table 1.

We employ several linear regression analysis methods to derive the best-fitting parameters of the \( \Gamma_{\text{tot}}-L/L_{\text{Edd}} \) correlation for the primary BASS sample. In all these fits, we assume a uniform uncertainty of 0.3 dex on \( L/L_{\text{Edd}} \) (following S08). The BCES(Y|X) method (Akritas & Bershady 1996) provides

\[ \Gamma_{\text{tot}} = (0.167 \pm 0.04) \log (L/L_{\text{Edd}}) + (2.00 \pm 0.05), \tag{2} \]

while the BCES bisector fit6 provides \( \alpha = 0.444 \pm 0.060 \) and \( \beta = 2.34 \pm 0.077 \). The FITEXY method, adapted to include intrinsic scatter (following Tremaine et al. 2002), provides \( \alpha = 0.167 \pm 0.029 \) and \( \beta = 2.004 \pm 0.038 \) (and an intrinsic scatter of 0.24) – in excellent agreement with the BCES(Y|X) result. Fig. 3 presents the BCES bisector and FITEXY best-fitting relations. Table 2 lists the best-fitting parameters for all three linear fits, as

3 Throughout this work, we define a correlation as ‘significant’ if the two-sided Spearman correlation test results in \( P < 0.1 \) per cent (corresponding to \( >3\sigma \)). Correlations with \( 0.1 < P < 1 \) per cent (i.e. \( \sim 2\sigma \)–\( 3\sigma \)) are referred to as ‘marginally significant’, in order to avoid a situation where

4 For the R09 study, we list the relation which relies on the ‘total’ sample, despite the fact that for \( \sim17 \) per cent of those AGN have \( C_{\lambda 5100} > 1549 \)-based determinations of \( M_{\text{BH}} \) (and therefore, \( L/L_{\text{Edd}} \)), which are known to be problematic (see Trakhtenbrot & Netzer 2012, and references therein). The relation derived in R09 for AGN with Hβ-based determinations of \( M_{\text{BH}} \) is much steeper, with \((\alpha, \beta) = (0.58, 2.57)\).

5 We stress that this value of the correlation coefficient \( r_s \) should not be directly compared with the slopes of the \( \Gamma_1-L/L_{\text{Edd}} \) relations reported by the aforementioned studies, despite their similarity.

6 All our BCES fits used 1000 realizations of the relevant data sets.
The grey diagonal lines represent the best-fitting residuals (i.e. $\sigma_P < 0.1$ per cent). We also tabulate the standard deviation of the resulting residuals show significant trends with $x$.

Moreover, these linear relations fail to reduce the considerable amount of scatter in the $\Gamma_{\text{tot}}-L/L_{\text{Edd}}$ plane: the standard deviations of the residuals, roughly $\sigma(\Delta) \gtrsim 0.25$, are comparable to the general standard deviation of $\Gamma_{\text{tot}}$ in our sample ($\sigma[\Gamma_{\text{tot}}] = 0.27$). Thus, there is little evidence that these linear relations provide a preferred description of the $\Gamma_{\text{tot}}-L/L_{\text{Edd}}$ parameter space, and/or the range in $\Gamma_{\text{tot}}$ seen in the BASS sample.

Despite the statistically significant (though weak) correlation found for the primary BASS sample as a whole, a closer inspection of the three different $M_{\text{BH}}$ subsets provides very limited evidence for such correlations within these subsets. In particular, the subsets with ‘direct’, ‘single-epoch’, and ‘$\sigma$’ determinations of $M_{\text{BH}}$ result in $P$-values of 94.6 per cent, 0.36 per cent, and 39.2 per cent, respectively (all based on Spearman correlation tests; see Table 1). We highlight the lack of a statistically significant correlation among the most reliable $M_{\text{BH}}$ determinations (i.e. the ‘direct’ subset) and the weak evidence for a correlation among the single-epoch subset, which most closely resembles the $M_{\text{BH}}$ estimation methodology of the aforementioned reference studies.

These apparently qualitatively inconsistent results – for the BASS sample as a whole and for the different $M_{\text{BH}}$ subsets – suggest much weaker dependence of $\Gamma_x$ on $L/L_{\text{Edd}}$, compared to those reported in previous studies (i.e. $\alpha \approx 0.16$ versus $\sim 0.31$).

We next examine the alternative determinations of $\Gamma_x$ and/or $L/L_{\text{Edd}}$ available for our sample, to further test whether we can establish any
(stronger) relations between these two quantities. In particular, we have examined relations between the \( L_{2-10} \)-based estimates of \( L/L_{\text{Edd}} \) (and \( f_{\text{bol}} = 0.10 \) per cent), and either \( \Gamma_{\text{int}} \), \( \Gamma_{\text{BAT}} \), or \( \Gamma_{\text{Edd}} \) – shown in Fig. 5. We have also used the alternative set of \( L/L_{\text{Edd}} \) estimates, in which \( L_{\text{bol}} \) is estimated from \( L_{2-10} \) and the bolometric corrections of either Marconi et al. (2004), or those of Vasudevan & Fabian (2007) – presented in the top two panels of Fig. 6. The \( L_{\text{BAT}} \)-based estimates of \( L/L_{\text{Edd}} \) (i.e. \( L_{\text{bol}} = 8.5 L_{\text{BAT}} \)) are presented in the bottom-left panel of Fig. 6, while the bottom-right panel presents the \( L_{5100} \)-based estimates of \( L/L_{\text{Edd}} \), for the subset of 128 AGN for which \( M_{\text{BH}} \) is determined from single-epoch spectroscopy of the broad H\( \beta \) line.

The results of all these tests are qualitatively similar to our main analysis of \( \Gamma_{\text{int}} \) versus \( L/L_{\text{Edd}}(L_{2-10}) \): large scatter, statistically significant correlations between \( \Gamma_{\text{int}} \) and \( L/L_{\text{Edd}} \) for the overall primary BASS sample (i.e. 228 AGN), but no correlation within any of the three \( M_{\text{BH}} \) subsets – as can be seen in the results of the formal correlation analyses (listed in Table 1). We particularly note that in all the cases we examined (i.e. all \( L/L_{\text{Edd}} \)), the most reliable ‘direct’ \( M_{\text{BH}} \) subset did not result in statistically significant correlations. The ‘single-epoch’ subset shows somewhat stronger evidence for correlations, with \( P \)-values \( \gtrsim 1 \) per cent in all cases, and a statistically significant (but weak) correlation for the case where \( \Gamma_{\text{Edd}} \) is considered \((P \lesssim 10^{-2} \text{ per cent}, r_c = 0.31)\); the best-fitting (\( \text{FITEXY} \)) relation has \( \alpha = 0.16 \). Another noteworthy exception is the lack of correlation between \( \Gamma_{\text{int}} \) and the \( L_{\text{BAT}} \)-based determinations of \( L/L_{\text{Edd}} \), even among the entire primary BASS sample (bottom panel of Fig. 3). Importantly, we find that the correlation between \( \Gamma_{\text{int}} \) and the \( L_{1000} \)-based estimates of \( L/L_{\text{Edd}} \), for the subset of AGN with single-epoch, broad H\( \beta \) determinations of \( M_{\text{BH}} \), is neither truly statistically significant \((P = 0.11 \text{ per cent})\) nor strong \((r_c = 0.29)\). We will revisit this subset when comparing our results with previous studies of the \( \Gamma_{\text{int}}-L/L_{\text{Edd}} \) relation (see Section 4.2).

We finally note that the BASS sample provides no compelling evidence for an ‘inversion’ of the \( \Gamma_{\text{int}}-L/L_{\text{Edd}} \) relation for low-\( L/L_{\text{Edd}} \).
Figure 5. Same as Fig. 3, but with alternative determination of the photon index, \(\Gamma_x\). Top left: \(\Gamma_{0.3-10}\) obtained from the full X-ray spectral model, fitted over the energy range 0.3–10 keV. Top right: \(\Gamma_{\text{BAT}}\) obtained from a power-law spectral model fitted over the Swift/BAT energy range of 14–195 keV. Bottom: \(\Gamma_{\text{eff}}\) obtained from a modified spectral model that ignores the high-energy exponential cut-off.

3.3 Additional tests for subsets of AGN

Finally, we examined several subsets of sources within our BASS sample, verifying that none of the choices we made in defining our sample, or our treatment of certain physically motivated spectral components, would have a significant effect on our conclusion. In particular, we tested for the existence of \(\Gamma_x-L/L_{\text{Edd}}\) correlations among: AGN with \(0.05 < z < 0.5\) – minimizing aperture effects; AGN with high-quality (SDSS) optical spectra; AGN with \(0.01 < L/L_{\text{Edd}} < 1\); AGN with no heavy obscuration \((\log (N_{\text{H}}/\text{cm}^{-2}) < 23)\); and AGN without warm absorbers. These subsets are described in Appendix A, and the results of the correlation tests are tabulated in Table A1. The qualitative results of this analysis are consistent with what we find for the primary BASS sample: for each subset, we find either no correlation, or alternatively, a weak correlation for all the AGN in that subset, while finding no correlations among sources with differing \(M_{\text{BH}}\) determination methods.

3.4 Relations between \(\Gamma_x\) and other AGN properties

We looked for relations between \(\Gamma_x\) and other key properties of the accreting SMBHs in our sample. Fig. 7 presents \(\Gamma_{\text{tot}}\) versus \(L_{2-10}, FWHM(H_\beta)\) (or \(H_\alpha\)), and \(M_{\text{BH}}\). The \(P\)-values associated with these correlation tests are listed in Table 1. None of these relations resulted in a statistically significant correlation. A qualitatively similar result was obtained when testing for correlations involving \(\Gamma_{0.3-10}\) or \(\Gamma_{\text{eff}}\). These results are in agreement with the findings of previous studies that investigated possible links between \(\Gamma_x\) and other AGN properties.

The broad dynamical range in \(L_{2-10}, L_{\text{bol}}, M_{\text{BH}},\) and \(L/L_{\text{Edd}}\) covered by our sample allows us to further investigate whether the mutual dependence between (some of) these quantities has any effects on the \(\Gamma_x-L/L_{\text{Edd}}\). To this end, we examined subsets of our sample for which one of these properties is controlled. Considering only the AGN with \(\log (L_{2-10}/\text{erg s}^{-1}) = 43.25-43.75\) (i.e. a bin of \(\pm 0.25\) dex around the median luminosity, with 71 sources), we find no evidence for a significant \(\Gamma_x-L/L_{\text{Edd}}\) correlation \((P = 2\text{ per cent for }\Gamma_{\text{tot}}, \text{and } >0.1\text{ per cent for all other cases})\). This should be compared to the highly significant correlations found when considering the entire luminosity range \((P \ll 10^{-3}\text{ per cent in all cases};\text{ see Table 1})\). A similar analysis for AGN with

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7 For the purposes of the test with the FWHM of broad Balmer lines, we focused only on the 149 AGN with single-epoch determinations of \(M_{\text{BH}}\) (i.e. ignoring the 50 AGN with ‘direct’ mass measurements).
log \((M_{\text{BH}}/M_{\odot}) = 7.75-8.25\) (again within ±0.25 dex of the median value; 64 sources) provides a qualitatively different result: the statistically significant \(\Gamma_x\)-\(L_{Edd}\) correlation holds for most cases \((P \approx 0.1\%\) for \(\Gamma_{\text{tot}}\) and \(\Gamma_{0.3-10} < 10^{-5}\%\) for \(\Gamma_{\text{bol}}\) and \(\Gamma_{\text{BAT}}\)). We stress that these two special subsets of AGN cover the same range in both \(L_{Edd}/L_{Edd}\) and \(\Gamma_x\) as does our primary BASS sample. This is only possible thanks to the broad range of \(L/L_{Edd}\) and \(M_{\text{BH}}\) provided through the BASS project (see K17).

These results, together with the fact that \(L_{Edd}/L_{Edd}\) is strongly correlated with \(L_{2.10}\) in our sample \((P \approx 10^{-5}\%\) per \(r_x = 0.29\)), suggest that the \(\Gamma_{\text{tot}}-L_{Edd}\) relation for the primary BASS sample may be at least partially driven by the trend with source luminosity.

4 DISCUSSION

4.1 The BASS \(\Gamma_x - L_{Edd}\) plane for different classes of AGN

Our analysis shows no evidence for a robust \(\Gamma_x-L_{Edd}\) relation among the subsets of AGN for which reliable estimates of \(M_{\text{BH}}\) (and therefore, of \(L_{Edd}\)) are available, while also showing evidence for a significant correlation among the BASS sample as a whole, as well as (marginal) evidence for a correlation among the broad-line sources. How could these qualitatively contradicting results be reconciled?

The study of Winter et al. (2012) has identified a similar discrepancy, when finding a strong \(\Gamma_x-L_{Edd}\) correlation only among the broad-line Swift/BAT-selected AGN in their sample. The interpretation put forward by that study suggested that the lower luminosity and/or lower \(L_{Edd}\)-absorbed AGN are found in a different accretion state. For our BASS sample, a closer inspection of Figs 1 and 3, suggests that the \(\sigma_x\) subset (i.e. AGN with no broad Balmer lines, and no direct \(M_{\text{BH}}\) determination) exhibits somewhat lower \(\Gamma_x\), compared with the other two \(M_{\text{BH}}\) subsets (see also Vasudevan, Mushotzky & Gandhi 2013). In addition, the studies of Fabian, Vasudevan & Gandhi (2008) and Fabian et al. (2009) showed that such narrow-line sources are predominantly low-\(L_{Edd}\) systems. The combined effect of these two trends is that the \(\sigma_x\) subset mainly extends towards the low-\(L_{Edd}\) end of the parameter space, which in turn results in statistically significant \(\Gamma_x-L_{Edd}\) correlations once this subset is included in the analysis.

Are these two trends driven by physical processes or by observational limitations (i.e. selection effects)? As suggested by Fabian et al. (2008), the tendency of obscured (narrow line) AGN towards low \(L_{Edd}\) is likely driven by the limited radiation pressure that low-\(L_{Edd}\) AGN exert on the surrounding dusty circumnuclear gas
when combined with the tendency of obscured AGN to have lower $L/L_{Edd}$ – produces the observed $\Gamma_x - L/L_{Edd}$ relation for our entire sample of BASS sources. One such scenario would be if obscured AGN have multiple partially covering (i.e. clumpy) absorption components (e.g. Cappi et al. 1996). In such a case, the measured $\Gamma_x$ might be flatter than the real underlying photon index. Thus, a scenario in which the lower $\Gamma_x$ of obscured sources is driven by physical effects beyond the $\Gamma_x - L/L_{Edd}$ correlation would require that many (or indeed, most) obscured sources would have (at least) two partially covering absorbing components. This is, arguably, a rather extreme scenario.

4.2 Comparison with previous studies

We demonstrated that our sample of 228 low-redshift, hard X-ray-selected AGN shows no significant evidence for a correlation between the hard X-ray photon index, $\Gamma_x$, and the normalized accretion rate, $L/L_{Edd}$, nor with other key AGN properties such as BH mass ($M_{BH}$) and/or hard X-ray luminosity ($L_{BAT}$). This stands in contrast to the findings of several studies. In what follows, we briefly summarize three such studies, which form the main reference for our comparison.

(i) Shemmer et al. (2008, S08) studied 35 high-luminosity, high-redshift quasars (at $z \sim 0.5$), for which the X-ray spectral analysis mostly relied on $XMM$-$Newton$ data in the observed-frame energy range 0.5–10 keV. BH masses were determined from broad H$\beta$ spectroscopy, using the same prescription we use here, and $L/L_{Edd}$ were calculated through the $L_{3.4}$-dependent prescription of Marconi et al. (2004), consistent with the $L_{3.4}$-dependent bolometric corrections we use here.

(ii) Risaliti et al. (2009, R09) analysed a sample of 343 moderate- to high-luminosity (43 $\lesssim \log \left[ L_x / \text{erg} \text{s}^{-1} \right] \lesssim 46.7$) SDSS quasars at $0.1 \lesssim z \lesssim 4.5$, with archival $XMM$-$Newton$ data (compiled by Young, Elvis & Risaliti 2009). The X-ray spectra, covering 0.5–10 keV, were fitted with an (absorbed) power-law model. BH masses were determined from either the H$\beta$, Mg $\equiv \lambda 2798$, or C iv $\lambda 1549$ broad emission lines (with 314 AGN having the more reliable H$\beta$- or Mg $\equiv$-based masses). Bolometric luminosities were derived by using a fixed-shape UV-optical SED, and a power-law X-ray SED (with $E_C = 100$ keV). The $\Gamma_x - L/L_{Edd}$ relations found for the subsets of AGN with either H$\beta$- or Mg $\equiv$-based $M_{BH}$ determinations are markedly different ($\alpha = 0.58$ and 0.24, respectively).

(iii) Brightman et al. (2013, B13) analysed a sample of 69 X-ray-selected, hard-line AGN from the $Chandra$ surveys in the E-CDF-S and COSMOS fields, covering 0.5 $\lesssim z \lesssim 2$ and 42.5 $\lesssim \log \left[ L_x / \text{erg} \text{s}^{-1} \right] \lesssim 45.5$. The sample was restricted to sources with more than 250 counts in their spectra. BH masses were obtained through either H$\alpha$- or Mg $\equiv$-based single-epoch estimators, which are generally consistent with those used here and in the other reference studies.

All these studies, which serve as primary reference studies for our work, focused on unobscured, broad-line AGN, for which $M_{BH}$ is determined through single-epoch spectroscopy of broad emission lines – comparable to our ‘single-epoch’ $M_{BH}$ subset. In addition, most of these studies employed a spectral model that includes only a single power law, with a minor absorption correction for a few sources.

The study of Winter et al. (2012) employed a more elaborate X-ray spectral model to a sample of broad-line Swift/BAT-selected AGN, and identified strong $\Gamma_x - L_X$ and $\Gamma_x - L/L_{Edd}$ relations, although the slope of the latter ($\alpha = 0.23$) is somewhat flatter than

(i.e. the dusty tori), which in turn result in increased levels of optical and X-ray obscuration. This issue is investigated in detail, in the context of the BASS sample and data, in a forthcoming study by Ricci et al. The outstanding question is therefore whether the somewhat lower $\Gamma_x$ seen in obscured AGN is driven by the $\Gamma_x - L/L_{Edd}$ correlation (the origin of which not yet well understood; see below), or rather by an unrelated physical and/or observational effect, which –

![Figure 7](image_url)

**Figure 7.** Testing for relations between X-ray photon index $\Gamma_{tot}$ and other key AGN properties. **Top:** $\Gamma_{tot}$ versus hard X-ray luminosity (14–150 keV) probed by Swift/BAT, $L_{BAT}$. **Centre:** $\Gamma_{tot}$ versus BH mass, $M_{BH}$. In both these panels, symbols are identical to Fig. 3. **Bottom:** $\Gamma_{tot}$ versus width of the broad Balmer lines (H$\beta$ or H$\alpha$) for those sources for which these data are available. No correlations are found between $\Gamma_{tot}$ and any of these properties.

This is illustrated in detail, in the context of the BASS sample and data, in a forthcoming study by Ricci et al. The outstanding question is therefore whether the somewhat lower $\Gamma_x$ seen in obscured AGN is driven by the $\Gamma_x - L/L_{Edd}$ correlation (the origin of which not yet well understood; see below), or rather by an unrelated physical and/or observational effect, which –
what is found for the optically selected quasars mentioned above. More recently, the study of Brightman et al. (2016) studied the $\Gamma_x - L/L_{\text{Edd}}$ relation in a sample of nine heavily obscured (mostly CT) AGN, for which precise $M_{\text{BH}}$ measurements are available from resolved megamaser kinematics. This analysis resulted in a significant correlation with best-fitting parameters that are consistent with those derived in the aforementioned studies of unobscured AGN.

Comparing these reference studies to our BASS analysis, we first note the higher quality and broader energy coverage of our BASS X-ray data. These allow for a much more elaborate and robust spectral decomposition, taking into account several physically motivated components, and provide a set of various determinations of the key quantities (i.e. $\Gamma_x$ and $L_{5100}$). We also note that our sample completely overlaps with the reference studies in terms of the range of $\Gamma_x$ and $L/L_{\text{Edd}}$ covered, and that it includes 174 broad-line AGN – the only class of AGN studied in the reference studies.

Although at face value our BASS analysis suggests a $\Gamma_x - L/L_{\text{Edd}}$ correlation which is similar to those found in the reference studies, we note two main differences. First, we stress that we find little evidence for any $\Gamma_x - L/L_{\text{Edd}}$ link among BASS sources for which $M_{\text{BH}}$ is determined from single-epoch spectra of broad emission lines – the only subset comparable with the reference studies. Even for this subset, the only statistically significant correlation we find is when using $\Gamma_{\text{Edd}}$ (which may be similar to the $\Gamma_x$ used in some of the reference studies). Moreover, the correlation involving $\Gamma_{\text{BAT}}$ – which could be thought of as comparable to what is measured for high-redshift sources (see S08) – is insignificant (although at $P = 0.16$ per cent). Second, the slopes of the best-fitting relations we derive for our entire BASS sample (Table 2) differ from those previously reported ($\alpha \approx 0.3$): we find $\alpha = 0.16$ for the $(\Gamma_x, L/L_{\text{Edd}})$ correlation analyses, but $\alpha \gtrsim 0.4$ for the BCES bivariate. The discrepancy between the different fitting methods probably reflects the large scatter in the $\Gamma_x - L/L_{\text{Edd}}$ plane.

To allow for a more direct comparison, we have derived yet another set of $\Gamma_x$ measurements which aims to resemble the analysis performed in previous studies. We re-fitted the X-ray data of 162 BASS AGN that have log ($N_H/cm^2$) $\leq 22$ with a simplified spectral model of an absorbed power law over the rest-frame energy range 2–10 keV. By ignoring any additional components (i.e. warm absorbers, reflection, Fe Kα), this model – and the chosen energy range – are similar to what was used in the aforementioned reference studies. We stress that these derived photon indices, $\Gamma_{\text{simple}}$, are not identical to $\Gamma_{0.3-10}$ (see Section 2.2), despite the similarity in the respective energy ranges, as $\Gamma_{0.3-10}$ was derived from a more elaborate spectral model. We further focus on those AGN for which $M_{\text{BH}}$ is determined through single-epoch spectroscopy of the broad Hβ emission line, and on the $L_{5100}$-based estimates of $L/L_{\text{Edd}}$.

Fig. 8 (top panel) plots these simplified photon indices ($\Gamma_{\text{simple}}$) against the $L_{5100}$-based estimates of $L/L_{\text{Edd}}$ for the relevant 119 AGN in our sample. In this case, we find a statistically significant ($P \approx 2 \times 10^{-4}$ per cent) yet, again, weak ($r_x = 0.383$) correlation between these two particular quantities. We recall that a similar analysis, with $L_{5100}$-based estimates of $L/L_{\text{Edd}}$ for the single-epoch, broad-Hβ subset, but with $\Gamma_{\text{tot}}$, yielded only a marginally significant correlation ($P = 0.1$ per cent; see Section 3.2 and Fig. 6). A formal correlation analysis results in ($\alpha, \beta$) = (0.906 ± 0.11, 2.79 ± 0.13), (0.304 ± 0.09, 2.00 ± 0.13), and (0.326 ± 0.08, 2.03 ± 0.10), for the BCES bivariate, BCES(Y[X], and FITEXY methods, respectively (with an intrinsic scatter of 0.2 added in the latter case). The best-fitting slopes of the latter two (Y[X]) relations are in excellent agreement with those reported by the main three reference studies. We stress that we find no significant correlation between $\Gamma_{\text{simple}}$ and the primary, $L_{2-10}$-based estimates of $L/L_{\text{Edd}}$ ($P = 0.8$ per cent), as seen in the bottom panel of Fig. 8. This is an important point, as some of the studies that reported strong $\Gamma_x - L/L_{\text{Edd}}$ correlations (e.g. R09, B13) relied, at least partially, on $L_{2-10}$-based determinations of $L/L_{\text{Edd}}$, and not on purely $L_{5100}$-based ones.

Thus, it appears that the photon index derived from a simplified X-ray spectral model of a power-law fit to the softer X-ray data of a subset of 119 unobscured (log ($N_H/cm^2$) $\leq 22$), broad-line AGN where $M_{\text{BH}}$ is determined from broad Hβ. Top: $L/L_{\text{Edd}}$ is estimated from $L_{5100}$. These data exhibit a statistically significant correlation, unlike what we found when considering $\Gamma_{\text{tot}}$ (cf. the bottom-right panel of Fig. 6). Bottom: $L/L_{\text{Edd}}$ is estimated from $L_{2-10}$. No correlation is found in this case.

### 4.3 Using $\Gamma_x$ as a BH growth indicator

As first pointed out by Shemmer et al. (2008), one of the exciting implications of a strong and tight relation between $\Gamma_x$ and $L/L_{\text{Edd}}$ is the possibility to use large X-ray surveys to construct nearly complete distributions of $L/L_{\text{Edd}}$, particularly for high-redshift sources.
in deep extragalactic fields, where this key quantity is otherwise hard to measure (e.g. Trakhtenbrot & Netzer 2012; Trakhtenbrot et al. 2016). The study of Brightman et al. (2013) can be considered as a demonstration of such an approach within a dedicated survey (COSMOS). Moreover, the recent study of Brightman et al. (2016) suggested that this approach may also be applicable to heavily obscured (CT) AGN, potentially providing a unique probe of the accretion rates among these elusive objects.

However, our sample and analysis highlight the limitations associated with using \( \Gamma_x \) measurements to predict \( L/L_{\text{Edd}} \). We first recall that the overall scatter in the \( \Gamma_x - L/L_{\text{Edd}} \) plane is large (~0.3 dex; see Figs 2–6), and that the few statistically significant relations we find between \( L/L_{\text{Edd}} \) and \( \Gamma_x \) are further demonstrated by the corresponding correlation analysis. The best-fitting BCES relation we find for our primary sample (i.e. 228 sources) is

\[
L/L_{\text{Edd}} = (0.71 \pm 0.27) \Gamma_{\text{tot}} - (2.44 \pm 0.48),
\]

which is consistent, within the considerable uncertainties, to the relation reported by S08 (their eq. 2). The uncertainties on the best-fitting parameters in eq. (3) are so large that for a given \( \Gamma_{\text{tot}} \) with zero measurement uncertainty, they predict values of \( L/L_{\text{Edd}} \) with a 1σ interpercentile range of 1.37 dex (i.e. the 16–84 per cent percentile range). Moreover, the corresponding FITEXY \( (L/L_{\text{Edd}} - \Gamma_{\text{tot}}) \) analysis suggests that a satisfactory fit, with \( \chi^2/\nu \approx 1 \), can only be obtained with the addition of a significant level of intrinsic scatter, exceeding 0.6.

Notwithstanding these limitations, it might still be possible to identify subsets of extremely high- or low-\( L/L_{\text{Edd}} \) AGN, probed by correspondingly extreme \( \Gamma_x \) (i.e. \( \Gamma_x \gtrsim 2.3 \) or \( \lesssim 1.2 \)). This is supported by the relatively clear separation between the peaks (and medians) of the distributions seen for \( \Gamma_x = 1.5 \) and 2.1 in Fig. 9.

We note, however, that such extreme \( \Gamma_x \) are only observed among
a minority of AGN, out to $z \sim 4$ (e.g. Just et al. 2007; Brandt & Alexander 2015; Cappelluti et al. 2016; Marchesi et al. 2016).

We conclude that the large scatter and weak correlations (at best) in the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ plane significantly hinder the prospects of using $\Gamma_{\text{x}}$ measurements to establish the distribution of $L_{\text{Edd}}$ among samples of high-redshift AGN.

4.4 Possible physical links between $\Gamma_{\text{x}}$ and $L_{\text{Edd}}$

Previous studies have tried to explain the positive $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ correlation through a picture where the increasing $L_{\text{Edd}}$ is causing an increased UV radiation, which in turn causes more efficient cooling in the corona. In principle, one may expect a similar trend of increasing $\Gamma_{\text{x}}$ with decreasing $M_{\text{BH}}$, as in the framework of geometrically thin, radiatively thick accretion discs this is also expected to increase the UV incident radiation (e.g. Davis & Hubeny 2006; Done et al. 2012; Davis & Laor 2011).

We however recall that our analysis showed no correlation between $\Gamma_{\text{x}}$ and $M_{\text{BH}}$ (Fig. 7). One way to accommodate this lack of trend with the aforementioned physical picture is if the X-ray-emitting corona is located closer to the disc for lower $M_{\text{BH}}$ systems, therefore reducing the amount of incident UV radiation. Such trends are indeed suggested by some reverberation mapping studies (see e.g. De Marco et al. 2013; Kara et al. 2013, and the review by Uttley et al. 2014).

We conclude that any scenario that connects the observed $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ relation to variations in the UV radiation field that is upscattered by the hot, X-ray emitting corona, should also account for the lack of observed relation between $\Gamma_{\text{x}}$ and $M_{\text{BH}}$ (and for that matter, with $L_{\text{X}}$; see again Fig. 7).

5 CONCLUSIONS

We presented a detailed analysis of the links between the hard X-ray photon index, $\Gamma_{\text{x}}$, and the (normalized) accretion rate, $L_{\text{Edd}}$, for a large sample of hard X-ray-selected, low-redshift AGN, as part of the BASS project. Our analysis was motivated by several earlier studies that identified significant, positive correlations between $\Gamma_{\text{x}}$ and $L_{\text{Edd}}$, over a broad range of redshifts. The low-redshift BASS sample allowed us to study these relations over a wide range of $L_{\text{AGN}}$–$M_{\text{BH}}$, and $L_{\text{Edd}}$. The high-quality and broad spectral coverage of the BASS data – unprecedented among studies that address the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ relation – allowed us to examine, for the first time in this context, the role of alternative determinations of the key quantities, and of the different methods used to derive them. Our main conclusions are as follows:

(i) Despite a significant amount of scatter, we find a weak (but statistically significant) correlation between $\Gamma_{\text{x}}$ and $L_{\text{Edd}}$ among our primary sample of 228 AGN. This correlation is robust to the choice of $\Gamma_{\text{x}}$.

(ii) The best-fitting $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ relations we obtain have flatter slopes than those reported by previous studies. Moreover, these best-fitting relations fail to reduce the scatter in the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ plane.

(iii) We find either no, or weak evidence for a $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ correlation when considering, separately, the subsets of AGN that differ in the method used to derive $M_{\text{BH}}$ (and therefore, $L_{\text{Edd}}$). In particular, we find no correlation for the subset of AGN with the most reliable, ‘direct’ mass estimates.

(iv) We find no statistically significant correlations between $\Gamma_{\text{x}}$ and either the $L_{\text{Edd}}$ estimates based on $L_{\text{Edd}}$, nor with $L_{2-10}$, $M_{\text{BH}}$, or the width of the broad Balmer emission lines.

(v) A $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ correlation that is consistent with those reported in previous studies does emerge, for a subset of broad-line AGN, when adopting a simplified, power law only spectral model fit to the lower energy X-ray data, and only when coupled with $L_{\text{Edd}}$ determinations that are based on the optical continuum emission.

(vi) We caution that the prospects of using the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ relation for deriving distributions of $L_{\text{Edd}}$ (and indeed $M_{\text{BH}}$) from deep X-ray surveys are limited due to the large scatter in the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ plane, the weakness of the correlations we find, and their dependence on specific methodological choices (i.e. bolometric corrections and X-ray energy ranges).

Our analysis clearly demonstrates the complexity of the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ plane, even for a uniformly selected sample of nearby AGN, with a rich collection of multiwavelength data, and a careful, elaborate spectral analysis. It appears that the previously reported strong relations between $\Gamma_{\text{x}}$ and $L_{\text{Edd}}$ may be, at least partially, driven by methodological choices (i.e. $L_{\text{Edd}}$ prescriptions) and/or limited spectral coverage and modelling in the X-ray regime. Our results hint that an underlying physical mechanism that links the shape of the X-ray SED with $L_{\text{Edd}}$ may indeed be at work, but is not yet well understood.

The existence and robustness of the $\Gamma_{\text{x}}$–$L_{\text{Edd}}$ relation may be re-evaluated with yet larger, unbiased samples of high X-ray-selected AGN, provided by ongoing surveys using the Swift and NuSTAR missions. In particular, the NuSTAR mission is providing high sensitivity and high spatial resolution hard X-ray data for hundreds of AGN ( Civano et al. 2015; Lansbury et al. 2017), reaching lower luminosities and/or higher redshifts than previous hard X-ray studies.

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APPENDIX A: RESULTS OF ADDITIONAL CORRELATION ANALYSIS

As noted in Section 3.3, we have performed a series of correlation tests for different subsets of AGN in order to verify that our main results are not driven by the particular choices made through the sample definition and spectral analysis parts of our work. Table A1 presents the results of the correlation hypothesis tests for these subsets, which include:

(i) AGN at 0.05 < z < 0.5 – a subset where the effects of spectroscopic aperture (relevant for $M_{BH}$, and there $L/L_{Edd}$ determination) are minimal. For this sample, we find no statistically significant correlation, neither when considering all $M_{BH}$ subsets (84 sources), nor when considering each of these subsets separately (see $P$-values in Table A1).

(ii) AGN with SDSS-based optical spectroscopy – a subset where the (relative and absolute) flux calibration is optimal, and where aperture effects are small and well understood. For this sample of 47 sources, we find a result similar to the general one: only the entire BH subset (i.e. single-epoch estimates based on SDSS spectra) does not show a correlation.

(iii) AGN with 0.01 < $L/L_{Edd}$ < 1 – a subset dominated by broad-line sources, which could in principle capture an underlying, positive $\Gamma_x - L/L_{Edd}$ correlation even if this relation flattens (or becomes an anticorrelation) at very low or very high $L/L_{Edd}$. For this sample of 195 sources (mostly broad-line AGN), we find a result similar to the general one: the entire sample results in a statistically significant, but weak correlation; the three $M_{BH}$ subsets (within the $L/L_{Edd}$ > 0.01 sample) do not show a correlation.
Table A1. BASS $\Gamma_x$-$L/L_{\text{Edd}}$ correlations: significance tests for minor subsets. Statistically significant $P$-values are highlighted in bold.

| $L_{\text{bol}}$ tracer | $\Gamma$ | Sub-sample | $N$ | $P$-value ( per cent) | $r_s$ |
|-------------------------|----------|------------|-----|-----------------------|------|
| $C \cdot L_{2-10}$ | $\Gamma_{\text{tot}}$ | $0.05 < z < 0.5$ – all | 84 | 8.1 | – |
|                        |          | – direct $M_{\text{BH}}$ | 9 | 61.3 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 64 | 9.9 | – |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 11 | 32.7 | – |
| $C \cdot L_{2-10}$ | $\Gamma_{\text{tot}}$ | SDSS spectra – all | 47 | 1.7 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 33 | 54.7 | – |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 14 | 10.4 | – |
| $C \cdot L_{2-10}$ | $\Gamma_{\text{tot}}$ | $0.01 < L/L_{\text{Edd}} < 1$ – all | 195 | $4.1 \times 10^{-2}$ | 0.251 |
|                        |          | – direct $M_{\text{BH}}$ | 29 | 63.2 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 140 | 0.83 | – |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 26 | 18.7 | – |
| $C \cdot L_{2-10}$ | $\Gamma_{\text{tot}}$ | $\log N_H < 23$ – all | 188 | $1.7 \times 10^{-3}$ | 0.308 |
|                        |          | – direct $M_{\text{BH}}$ | 26 | 65.5 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 137 | 0.25 | – |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 25 | 39.5 | – |
| $C \cdot L_{2-10}$ | $\Gamma_{0.3-10}$ | $\log N_H < 23$ – all | 188 | $1.6 \times 10^{-5}$ | 0.371 |
|                        |          | – direct $M_{\text{BH}}$ | 26 | 91.4 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 137 | 0.067 | 0.287 |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 25 | 39.5 | – |
| $C \cdot L_{2-10}$ | $\Gamma_{\text{nEc}}$ | $\log N_H < 23$ – all | 188 | $5.3 \times 10^{-6}$ | 0.384 |
|                        |          | – direct $M_{\text{BH}}$ | 26 | 93.8 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 137 | 0.0058 | 0.337 |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 25 | 88.0 | – |
| $C \cdot L_{2-10}$ | $\Gamma_{\text{tot}}$ | No warm absorbers – all | 184 | $3.4 \times 10^{-3}$ | 0.301 |
|                        |          | – direct $M_{\text{BH}}$ | 20 | 69.1 | – |
|                        |          | – single-epoch $M_{\text{BH}}$ | 115 | 0.65 | – |
|                        |          | – $\sigma_*$-based $M_{\text{BH}}$ | 49 | 39.2 | – |

(iv) AGN with $\log (N_H/cm^{-2}) < 23$ – a subset where the effects of Compton scattering on the X-ray spectral decomposition are minimal. For this sample of 188 sources we find, again, a result similar to the general one: only the entire sample results in a significant correlation, while the two main $M_{\text{BH}}$ subsets (within the $\log (N_H/cm^{-2}) < 23$ sample) do not show a correlation. For this subset, we also tested correlations involving $\Gamma_{0.3-10}$ and $\Gamma_{\text{nEc}}$, which are expected to be most sensitive to a Compton scattering component. Indeed, we find that for this subset the correlations involving $\Gamma_{0.3-10}$ and $\Gamma_{\text{nEc}}$ are somewhat stronger that those found with $\Gamma_{\text{tot}}$. However, the qualitative outcome remains identical.

(v) AGN without warm absorbers – the presence of significant ionized absorption in the X-ray spectrum might lead to deviation of $\Gamma_x$ from the intrinsic value. For this sample of 184 sources, we again find a result consistent with the general one: the entire sample results in a significant correlation, while the two main $M_{\text{BH}}$ subsets do not show a correlation.

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