OPTIMUM BASETHICKNESS DETERMINATION OF A BACKILLUMINATED SILICON SOLAR CELL: IRRADIATION EFFECT

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Abstract

In this paper, we propose a method to determine the optimum thickness of a bifacial silicon solar cell back illuminated and under irradiation. The expressions of back surface recombination velocity of excess minority carrier depending on both irradiation energy (ϕp) and damage coefficient (kl) are established. From their plots, base optimum thickness is deduced, and its relationships with both irradiation energy (ϕp) and damage coefficient (kl) are modeled.

Introduction:

The damage imposed on a solar cell, by the irradiation flux of charged particles (ϕp), in certain experimental laboratory conditions or in space, is modeled through the expressions of diffusion coefficient or lifetime [1] of the excess minority carriers in the semiconductor material.

Our study concerns the silicon solar cell (n+/p/p+/n+) previously subjected to various doses of radiation and illuminated by the back side with polychromatic light [2].

The density of the excess minority carrier is obtained by solving the diffusion equation in the base (p) of the solar cell, provided with the boundary conditions at the junction surface (p+/p) and on the back surface (p+). These boundary conditions bring to the excess minority carriers recombination velocity, respectively (Sf) at the junction and (Sb) on the rear [3] [4] [5] [6].

The photocurrent is calculated and obtained as a function of these recombination velocities (Sf, Sb) and the parameters (ϕp, kl) of the irradiation imposed on the solar cell of (H) thick base.

The study of the photocurrent as a function of the (Sf) recombination velocity at the junction, for each irradiation parameter yields to extract the expressions of the (Sb) recombination velocity on the rear face [7]. The analysis of these recombination velocity expressions through their graphic representations as a function of the solar cell thickness (H)[8] yields to obtain the optimum thickness under each irradiation parameter and mathematically modeled.

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Theory:

Minority carrier density:

Figure 1 shows an (n+/p/p+) silicon solar cell[5],[9], [10],back surface illuminated[11], [12](double side surface field) with polychromatic light and submitted to irradiation[13] [14] [15].

![Diagram of an (n+/p/p+) silicon solar cell](image)

**Figure 1:** n⁺-p-p⁺ solar cell back surface illuminated with polychromatic light and under irradiation.

The continuity equation governing the excess minority carrier's phenomena i.e. diffusion in the bulk and surfaces recombination, is defined by the following expression:

\[
D(kl, \phi_p) \frac{\partial^2 G_{ar}(x, kl, \phi_p)}{\partial x^2} - \frac{\delta_{ar}(x, kl, \phi_p)}{\tau} + G_{ar}(x) = 0
\]  

\(D\) represents the excess minority carrier diffusion. When the material is submitted to the external conditions, it becomes effective diffusion coefficient and it is then dependent of parameters such as, magnetic field[16], temperature[17], frequency[18] [19]. Intrinsic parameters as doping rate [5] also influence the diffusion coefficient as well as grain sizes and grain recombination velocity [20] [21] [22] [23].

The well-known Einstein's relation links the diffusion coefficient \(D(kl, \phi_p)\) [24] [25] to the lifetime \(\tau\) and the diffusion length \(L\), for given irradiation \(\phi_p\) and intensity \(kl\).

\[
D(kl, \phi_p) = \frac{L^2(kl, \phi_p)}{\tau} \quad (2)
\]

\[
L(kl, \phi) = \frac{1}{\left( \frac{1}{L_0^2} + kl \cdot \phi \right)^{\frac{3}{2}}} \quad (3)
\]

\(D(kl, \phi_p)\): is the diffusion coefficient of the electrons in the base under irradiation.

\(L(kl, \phi_p)\): is the diffusion length of the excess minority carrier in the base as a function of the irradiation energy and the damage coefficient.

\(\delta(x, kl, \phi)\): represents the density of excess minority carrier in the base resulting from polychromatic illumination and under the influence of irradiation.

\(G_{ar}(x)\) is the rate of generation of excess minority carrier at depth \(x\) in the base under polychromatic illumination. Its expression is given by[26]:

\[
G_{ar}(x) = n \times \sum_{i=1}^{3} a_i \times e^{-b_i(H-x)} \quad (4)
\]
x : is the depth in the base of the solar cell.

\( \delta(x, kl, \phi_p) \) is determined from the resolution of equation (1) and is given in the following form:

\[
\delta_{ar}(x, kl, \phi_p) = A_{ar} \times c h \left[ \frac{x}{L(kl, \phi_p)} \right] + B_{ar} \times s h \left[ \frac{x}{L(kl, \phi_p)} \right] + \sum_{i=1}^{3} K_i \times e^{-b_i(H-x)} \tag{5}
\]

With \( K_i = \frac{n \times [L(kl, \phi_p)]^3 \times a_i}{D(kl, \phi_p) \times b_i^2 \times L^2(kl, \phi_p) - 1} \) \tag{6}

The \( A_{ar} \) and \( B_{ar} \) coefficients are determined from the boundary conditions:

At the junction (\( x=0 \)) :

\[
D(kl, \phi_p) \frac{\partial \delta_{ar}(x, kl, \phi_p)}{\partial x} \bigg|_{x=0} = S_f \times \delta_{ar}(0, kl, \phi_p) \tag{7}
\]

\( S_f \) is the excess minoritycarrier’s recombination velocity at the junction. It characterizes the operating point of the solar cell varying from short-circuit to open circuit [3] [4] [23].

At the rear face (\( x=H \)) :

\[
D(kl, \phi_p) \frac{\partial \delta_{ar}(x, kl, \phi_p)}{\partial x} \bigg|_{x=H} = -S_b(kl, \phi_p) \times \delta_{ar}(H, kl, \phi_p) \tag{8}
\]

\( S_b \) is the back surface recombination velocity of excess minority carrier[3],[25],[27],[28] resulting from the electric field created by the (p/p+) high-low junction [11]. It characterizes the behavior of the excess minority carrierat the base back surface.

**Results And Discussions:**

**Influence of the irradiation on photocurrent density:**
The photocurrent density is given by the following expression:

\[
J_{ph_{ar}}(S_f, kl, \phi_p) = q \times D(kl, \phi_p) \times \left[ \frac{B_{ar}}{L(kl, \phi_p)} + \sum_{i=1}^{3} k_i \times b_i \times e^{-b_i \times H} \right] \tag{9}
\]

Figures 2 and 3 show the profiles of the photocurrent density as a function of the recombination velocity at the junction\( (S_f) \) for different energy flows and for different damage coefficients. Figure 4 represents the profile of the photocurrent density as a function of the recombination velocity at the junction for different thickness of the base values.
**Figure 2:** Photocurrent density as a function of the junction recombination velocity for different energy flows with $k_l=6\text{cm}^{-2}/\text{MeV}$, $H=170\mu\text{m}$.

**Figure 3:** Photocurrent density as a function of the junction recombination velocity for different damage coefficients with $\phi_p=100\text{MeV}$, $H=170\mu\text{m}$.
Figure 4: Photocurrent density as a function of the junction recombination velocity for different base thickness with $\phi_p=100\text{MeV}$, $kl=6\text{cm}^2/\text{MeV}$.

At Low $S_f$ value, the solar cell operates under open circuit condition. It is indicated on all the figures (2, 3, 4) that the photocurrent remained zero. Then when $S_f$ increases, the photocurrent increases until reaching a plateau, which corresponds to the short circuit current regardless of the flow of irradiation, its intensity or the thickness of the base. The short circuit current increases with both the intensity of irradiation and the base thickness and decreases with the irradiation flux.

The expression of $S_b$ is then deduced by deriving equation (9) of the photocurrent density whatever ($f_p$, $k_l$, and $H$) with respect to $S_f$ and at large $S_f$ ranges [3] [4].

$$\left[ \frac{\partial J_{ph}(S_f, k_l, \phi_p)}{\partial S_f} \right]_{S_f \phi \times 10^4} = 0$$

The resolution of this equation (10) thus gives respectively two expressions of the back surface recombination velocity as:

$$S_b1(H, k_l, \phi_p) = - \frac{D(k_l, \phi_p)}{L(k_l, \phi_p)} \times \tanh \left( \frac{H}{L(k_l, \phi_p)} \right)$$

$$S_b2(H, k_l, \phi_p) = \frac{D(k_l, \phi_p) \times b_j \times \left[ -1 + ch \left( \frac{H}{L(k_l, \phi_p)} \right) \times e^{-b_H} \right] + \frac{1}{L(k_l, \phi_p)} \times sh \left( \frac{H}{L(k_l, \phi_p)} \right) \times e^{-b_H}}{1 - L(k_l, \phi_p) \times \left[ b_j \times sh \left( \frac{H}{L(k_l, \phi_p)} \right) \times e^{-b_H} + ch \left( \frac{H}{L(k_l, \phi_p)} \right) \right]}$$

Influence of the irradiation on the optimum thickness ($H_{opt}$) by the technique of the back surface recombination velocity

Figures 5 and 7 below show the profiles of the recombination velocity as a function of the thickness of the base for different energy flows and for different damage coefficients.
Figure 5: Back surface recombination velocity as a function of the base thickness for different energy flows with $k_l=6\text{cm}^2/\text{MeV}$.

Table 1 summarizes the obtained optimum base width of the solar cell while varying the energy flows.

| $\Phi_p(\text{MeV})$ | 100  | 125  | 150  | 175  | 200  |
|----------------------|------|------|------|------|------|
| Hopt(cm)             | 0.0133 | 0.0129 | 0.0126 | 0.0124 | 0.0120 |
| Sb1(cm/s)            | 2221.4 | 2189.3 | 2160.7 | 2128.6 | 2101.8 |
| Sb2(cm/s)            | $1.6571\times10^3$ | $1.6314\times10^3$ | $1.6086\times10^3$ | $1.5829\times10^3$ | $1.5614\times10^3$ |

Tableau 1: Hopt values for different irradiation energy flows.

Figure 6 shows the plot of optimum base thickness versus irradiation energy flow.

![Optimum thickness versus irradiation energy](image)

The decrease of (Hopt) optimum thickness of the base with irradiation energy is represented by a linear function. The obtained modelling relation is given as follow:

$$Hopt(cm) = a \times \Phi_p + b \quad (13)$$

$$a = -10^{-5} \text{cm/MeV} \quad b = 0.0146 \text{cm}$$
Figure 7 is the representation of the technique of intercept of recombination velocity curves that allows the extraction of the abscissa, which is the optimum base thickness while varying the intensity of irradiation (kl).

![Figure 7](image)

**Figure 7:** Back surface recombination velocity as a function of the base thickness for different damage coefficient with ϕ_p=100 MeV.

Table 2 summarizes Hopt values based on variations in irradiation intensity (kl).

| kl (cm^2/MeV) | 6     | 8     | 10    | 11    | 12    |
|---------------|-------|-------|-------|-------|-------|
| Hopt (cm)     | 0.0134| 0.0129| 0.0124| 0.0122| 0.0119|
| Sb1 (cm/s)    | 2219.6| 2180.4| 2141.1| 2117.9| 2100  |
| Sb2 (cm/s)    | 1.6557x10^5| 1.6243x10^5| 1.5929x10^5| 1.5743x10^5| 1.5600x10^5|

![Tableau 2](image)

**Tableau 2:** Hopt values for different damage coefficient.

Figure 8: Optimum thickness versus damage coefficient.

The optimum thickness is modeled as a decreasing linear function depending on the irradiation damage coefficient. The relationship is as follows:

\[ H_{opt} = m \times kl + n \]  

\[ m = -2 \times 10^{-4} \text{ cm/MeV} \]  

\[ n = 0.0149 \text{ cm} \]
The damages caused by the irradiation of a solar cell then lead to a reduction of the optimum thickness necessary for the production of an important photocurrent.

Knowing that the excess minority carrier recombination velocity depends on diffusion parameters (D, L) and optical parameters (monochromatic or polychromatic absorption coefficient of the material), results on the determination of optimum thickness were produced. These results take into account external factors that influence diffusion parameters (L, D). These include temperature[29], magnetic field[30], irradiation flow by charged particles[14], frequency of modulation[31] of optical or electrical signal, or combination of these factors[32],[33]. The conditions of manufacture of the semiconductor material through the doping rate also affect the diffusion coefficient [8].

The absorption coefficient (variable or constant) influences the recombination velocity in the rear face, also allows to produce results on the optimum thickness [34].

The different types of solar cells including the horizontal or vertical junction [35],[36] have been studied. Future investigations will take into account the 3D study model (grain boundary recombination velocity and grain size) and the combination of several external factors.

Conclusion:-
In this paper, we have proposed a method for determining the optimal thickness of the bifacial solar cell back surface illuminated and under irradiation. The expressions of the excess minority carrier in the base and the photocurrent density have been proposed. Calibration curves of the photocurrent density were plotted versus the junction recombination rate for different values of the irradiation energy (ϕp), the base thickness H and the damage coefficient (kl). The expressions of the excess minority carrier recombination velocity at the back surface, have been deduced and resolved graphically, to obtain the optimum thickness at intercept points of the curves. The correlation between the irradiation energy, the damage coefficient and the optimal thickness of the solar cell has been established.

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