Quantum Mechanical Reflection and Transmission Coefficients for a Particle through a One-Dimensional Vertical Step Potential

Rohit Gupta, Tarun Singhal, Dinesh Verma

Abstract: In this paper, we illustrate an application of the Laplace transformation for finding the quantum mechanical Reflection and Transmission coefficients for a particle through a one-dimensional vertical step potential. Quantum mechanics is one of the branches of physics in which the physical problems are solved by algebraic and analytic methods. By applying the Laplace transformation, we can find the quantum mechanical Reflection and Transmission coefficients for a particle through a one-dimensional vertical step potential. Generally, the Laplace transformation has been applied in different areas of science and engineering and makes it easier to solve the problems in engineering applications. It is a mathematical tool which has been put to use for solving the differential equations without finding their general solutions. It has applications in nearly all science and engineering disciplines like analysis of electrical circuits, heat and mass transfer, fluid dynamics, nuclear physics, process controls, quantum mechanical problems, etc.

Index terms: Quantum mechanical Reflection and Transmission coefficients, one-dimensional vertical step potential, Laplace transformation.

I. INTRODUCTION

Consider a particle of mass m and total energy E which is moving from a region where the potential is zero to a region where the potential is constant, then the potential function is represented as $V(z) = 0$ for $z < 0$ and $V(z) = V_o$ for $z > 0$. It is clear that the potential function undergoes only one discontinuous change at $z = 0$ as shown in figure 1. The potential function of this type is known as a vertical step potential. Hence when the force field acting on a particle is zero everywhere except in the limited region, it is known as a potential step or single step barrier or vertical step potential [1-2]. If we consider a beam of particles (say electrons) of energy E moving from left to right, i.e. along the positive direction of Z-axis and impinging from the left on a vertical step potential at $z = 0$, then according to quantum mechanics, the particles behave like a wave moving from left to right and face a sudden shift in the potential at $z = 0$. The problem is analogous when light strikes a sheet of glass where there is a shift in the index of refraction and the wave is partially transmitted.

Hence in this problem, the electrons will be partially reflected and partially transmitted at the discontinuity.

In this paper, an application of the Laplace transformation is illustrated for finding the quantum mechanical Reflection and Transmission coefficients for a particle through a one-dimensional vertical step potential. The reflection coefficient $R$ is a measure of the fraction of electrons reflected at the potential discontinuity and is defined as the ratio of reflected probability flux to the incident probability flux. The transmission coefficient $T$ is a measure of the fraction of electrons transmitted through the potential discontinuity and is defined as the ratio of transmitted probability flux to the incident probability flux. The probability flux is defined as the product of velocity of the particle in the given region and the probability density of particle in that region [3-4].

II. DEFINITION OF LAPLACE TRANSFORMATION

The Laplace transformation [5] of $g(z)$, which is defined for real numbers $z \geq 0$, is denoted by $L \{g(z)\}$ or $G(r)$ and is defined as $L \{g(z)\} = G(r) = \int_0^\infty e^{-rz} g(z) \, dz$, provided that the integral exists, i.e. convergent. If the integral is convergent for some value of $r$, then the Laplace transform of $g(z)$ exists otherwise not, where $r$ is the parameter which may be a real or complex number and $L$ is the Laplace transform operator.

1.1. Laplace Transformation of Elementary Functions

a. $L \{1\} = \frac{1}{r}, \quad r > 0$

b. $L \{z^n\} = \frac{n!}{r^{n+1}}, \quad r > 0$

where $n = 0, 1, 2, \ldots$

c. $L \{e^{cz}\} = \frac{1}{r - c}, \quad r > c$
Quantum Mechanical Reflection and Transmission Coefficients for a Particle through a One-Dimensional Vertical Step Potential

1.2. Laplace Transformation of Derivative of a function

If the function \( g(z) \), \( z \geq 0 \) is having an exponential order, that is if \( g(z) \) is a continuous function and is a piecewise continuous function on any interval, then the Laplace transform of derivative [6-7] of \( g(z) \) i.e. \( L \{ g'(z) \} \) is given by

\[
L \{ g'(z) \} = \left[ 0 - g(0) \right] - \int_0^\infty e^{-rz} g(z)dz,
\]

Or \( L \{ g'(z) \} = -g(0) + r \int_0^\infty e^{-rz} g(z)dz \)

Or \( L \{ g'(z) \} = rL \{ g(z) \} - g(0) \)

Or \( L \{ g'(z) \} = rg'(0) - g(0) \)

Since \( L \{ g(z) \} = rL \{ g(z) \} - g(0) \), therefore,

\[
L \{ g''(z) \} = rL \{ g'(z) \} - g(0)
\]

Or

\[
L \{ g''(z) \} = rL \{ g(z) \} - g(0) - g'(0)
\]

Or

\[
L \{ g''(z) \} = r^2L \{ g(z) \} - rg(0) - g'(0)
\]

Or

\[
L \{ g''(z) \} = r^2g(r) - rg(0) - g'(0)
\]

And so on.

1.3. Inverse Laplace Transformation

The inverse Laplace transform [6-7] of the function \( G(r) \) is denoted by \( L^{-1} \{ G(r) \} \) or \( g(z) \). If we write \( L \{ g(z) \} = G(r) \), then \( L^{-1} \{ G(r) \} = g(z) \), where \( L^{-1} \) is called the inverse Laplace transform operator.

1.3.1. Inverse Laplace Transformations of Some Functions

a. \( L^{-1} \{ \frac{1}{r} \} = 1 \)

b. \( L^{-1} \{ \frac{1}{(r-c)} \} = e^{cz} \)

c. \( L^{-1} \{ \frac{1}{r^2+c^2} \} = \frac{1}{c} \sin cz \)

d. \( L^{-1} \{ \frac{1}{r^2+cz^2} \} = \cos cz \)

e. \( L^{-1} \{ \frac{1}{r^2} \} = \frac{1}{(n-1)!} n > 0 \)

III. METHODOLOGY

The time-independent Schrödinger equation [8] in one-dimension is written as:

\[
D_z^2 \psi(z) + \frac{2m}{\hbar^2} [E - V(z)] \psi(z) = 0 \quad \text{...(1)}
\]

This equation is second order linear differential equation. In this equation \( \psi(z) \) is probability wave function of the particle incident from left on the one-dimensional vertical step potential of height \( V_0 \) (a positive constant) at \( z = 0 \), \( E \) is the total energy of the particle and \( V(z) \) is the potential energy of the particle.

If \( \psi_L(z) \) and \( \psi_R(z) \) are the probability wave functions to the left and the right side of the vertical step potential at \( z = 0 \), then these wave functions and their first order derivatives are continuous [4, 9] at \( z = 0 \) i.e.

\[
\psi_L(0) = \psi_R(0) = (\text{say})
\]

and \( D_z \psi_L(0) = D_z \psi_R(0) = B \text{ (say)} \), where \( A \) and \( B \) are constants.

In the region, \( z < 0 \), \( V(z) = 0 \), therefore, in this region the time-independent Schrödinger equation in one-dimensions is written as

\[
D_z^2 \psi_L(z) + \frac{2m}{\hbar^2} \psi_L(z) = 0 \quad \text{...(3)}
\]

Taking Laplace transforms of equation (3), we get

\[
r^2 \overline{\psi_L(r)} - r \psi_L(0) - D_z \psi_R(0) + \frac{2mE}{\hbar^2} \overline{\psi_L(r)} = 0.
\]

Or

\[
r^2 \overline{\psi_L(r)} - r A - B + \frac{2mE}{\hbar^2} \overline{\psi_L(r)} = 0.
\]

Rearranging the equation, we get

\[
\overline{\psi_L(r)} = \frac{r A + B}{r^2 + k_1^2} \quad k_1 = \frac{2mE}{\hbar^2}
\]

Or \( \overline{\psi_L(r)} = \frac{r A}{r^2 + k_1^2} + \frac{B}{r^2 + k_1^2} \)

Taking inverse Laplace transform, we get

\[
\psi_L(z) = A \cos k_1 z + \frac{B}{k_1} \sin k_1 z \quad \text{...(4)}
\]

Writing \( \cos k_1 z \) and \( \sin k_1 z \) in terms of exponentials, we get

\[
\psi_L(z) = \frac{A e^{ik_1 z} + e^{-ik_1 z}}{2} + \frac{B}{k_1} \left( e^{ik_1 z} - e^{-ik_1 z} \right)
\]

Or

\[
\psi_L(z) = \left( \frac{A}{2} + i \frac{B}{2k_1} \right) e^{ik_1 z} + \left( \frac{A}{2} - i \frac{B}{2k_1} \right) e^{-ik_1 z}
\]

Or

\[
\psi_L(z) = \frac{\left( \frac{e^{ik_1 z} + e^{-ik_1 z}}{2} \right)}{2} \left( \frac{e^{ik_1 z} - e^{-ik_1 z}}{2i} \right)
\]

And \( \frac{\left( \frac{e^{ik_1 z} + e^{-ik_1 z}}{2} \right)}{2} \left( \frac{e^{ik_1 z} - e^{-ik_1 z}}{2i} \right) \) represent the incident and the reflected waves in the region \( z < 0 \) i.e.

\[
\psi_{in}(z) = \left( \frac{A}{2} - i \frac{B}{2k_1} \right) e^{ik_1 z}
\]

and

\[
\psi_{re}(z) = \left( \frac{A}{2} + i \frac{B}{2k_1} \right) e^{-ik_1 z}
\]

Now, in the region, \( z > 0 \), depending on the values of \( E \) and \( V_0 \), two possibilities arise: either \( V(z) = V_0 < E \) or \( V(z) = V_0 > E \).

Case I: \( E > V_0 \)

In this case, in the region, the \( z > 0 \), \( V(z) = V_0 < E \), therefore, the time-independent Schrödinger equation in one-dimension is written as

\[
D_z^2 \psi_R(z) + k_2^2 \psi_R(z) = 0 \quad \text{...(6)}
\]

where \( k_2 = \sqrt{\frac{2m(E-V_0)}{\hbar^2}} \) is real.

Taking Laplace transforms of equation (3), we get

\[
r^2 \overline{\psi_R(r)} - r \psi_R(0) - D_z \psi_R(0) + k_2^2 \overline{\psi_R(r)} = 0.
\]

Or

\[
r^2 \overline{\psi_R(r)} - r A - B + k_2^2 \overline{\psi_R(r)} = 0.
\]

Rearranging the equation, we get

\[
\overline{\psi_R(r)} = \frac{r A + B}{r^2 + k_2^2} \quad k_2 = \sqrt{\frac{2m(E-V_0)}{\hbar^2}}
\]

Or

\[
\overline{\psi_R(r)} = \frac{r A}{r^2 + k_2^2} + \frac{B}{r^2 + k_2^2}
\]

Taking inverse Laplace transform, we get

\[
\psi_R(z) = A \cos k_2 z + \frac{B}{k_2} \sin k_2 z \quad \text{...(4)}
\]

\[
\psi_R(z) = \frac{A e^{ik_2 z} + e^{-ik_2 z}}{2} + \frac{B}{k_2} \left( e^{ik_2 z} - e^{-ik_2 z} \right)
\]

Or

\[
\psi_R(z) = \left( \frac{A}{2} + i \frac{B}{2k_2} \right) e^{ik_2 z} + \left( \frac{A}{2} - i \frac{B}{2k_2} \right) e^{-ik_2 z}
\]

Or

\[
\psi_R(z) = \frac{\left( \frac{e^{ik_2 z} + e^{-ik_2 z}}{2} \right)}{2} \left( \frac{e^{ik_2 z} - e^{-ik_2 z}}{2i} \right)
\]

And \( \frac{\left( \frac{e^{ik_2 z} + e^{-ik_2 z}}{2} \right)}{2} \left( \frac{e^{ik_2 z} - e^{-ik_2 z}}{2i} \right) \) represent the incident and the reflected waves in the region \( z > 0 \).
Writing \( \cos k_2z \) and \( \sin k_2z \) in terms of exponentials, we get
\[
\psi_R(z) = A \frac{e^{ik_2z} + e^{-ik_2z}}{2} + B \frac{e^{ik_2z} - e^{-ik_2z}}{2i} \]
Or
\[
\psi_R(z) = \left( \frac{A}{2} + \frac{B}{2i}k_2 \right) e^{ik_2z} + \left( \frac{A}{2} - \frac{B}{2i}k_2 \right) e^{-ik_2z} \]
Or
\[
\psi_L(z) = \left( \frac{A}{2} - \frac{B}{2i}k_2 \right) e^{ik_2z} + \left( \frac{A}{2} + \frac{B}{2i}k_2 \right) e^{-ik_2z} \ldots (7) \]
In equation (7), the first term on R.H.S.i.e. \( \left( \frac{A}{2} - \frac{B}{2i}k_2 \right) \) represents the transmitted wave in the region \( z > 0 \) i.e.
\[
\psi_{tr}(z) = \left( \frac{A}{2} - \frac{B}{2i}k_2 \right) e^{ik_2z} . \]
Since \( \left( \frac{A}{2} + i \frac{B}{2k_2} \right) \) represents the coefficient of a beam incident from right on the potential step, which is not physical, therefore, \( A + i \frac{B}{2k_2} = 0 \)
Or \( A = -i \frac{B}{2k_2} \)
Or \( B = ik_2A \ldots (8) \)
Using equations (8), we can write
\[
\psi_{in}(z) = \frac{A}{2} \left( 1 + \frac{k_2}{k_1} \right) e^{ik_1z} , \]
\[
\psi_{re}(z) = \frac{A}{2} \left( 1 - \frac{k_2}{k_1} \right) e^{-ik_1z} \]
and
\[
\psi_{te}(z) = Ae^{ik_2z} . \]
The quantum mechanical reflection coefficient \( R \) is given by\[^{[9]}\]
\[
R = \frac{\text{reflected probability flux}}{\text{incident probability flux}} \]
or \( R = \frac{\psi_{te} \overline{\psi_{re}}}{\psi_{in}}, \) where \( v_i \) is the velocity of particle in the region \( z < 0 \) and \( \overline{\psi_{re}} \) is the complex conjugate of \( \psi_{re} \) and \( \psi_{in} \) is the complex conjugate of \( \psi_{in} \).
\[
R = \frac{A \left( \frac{1}{2} \right) \left( \frac{1}{2} \frac{k_2}{k_1} \right) e^{ik_2z} - \frac{A}{2i} \left( \frac{1}{2} \frac{k_2}{k_1} \right) e^{-ik_2z}} {2 \left( \frac{1}{2} \frac{1}{2} \frac{k_2}{k_1} \right) e^{ik_1z} - 2 \frac{1}{2} \frac{1}{2} \frac{k_2}{k_1} e^{-ik_1z}} \]
Or \( R = \frac{\left( 1 - \frac{k_2}{k_1} \right) \left( 1 + \frac{k_2}{k_1} \right) e^{-ik_1z}} {1 + \frac{k_2}{k_1} \left( 1 - \frac{k_2}{k_1} \right) e^{ik_1z}} \)
On simplifying, we get
\[
R = \frac{\left| k_1 - k_2 \right|^2} {\left| k_1 + k_2 \right|^2} \ldots \ldots \ldots \ldots \ldots (9) \]
It is clear that the value of \( R \) is positive and less than one. The equation (9) provides the expression for the quantum mechanical reflection coefficient at the vertical step potential with \( E > V_0 \).

The quantum mechanical transmission coefficient \( T \) is given by\[^{[9]}\]
\[
T = \frac{\text{transmitted probability flux}}{\text{incident probability flux}} \]
or \( T = \frac{\psi_{te} \overline{\psi_{tr}}}{\psi_{in} \overline{\psi_{in}}}, \) where \( v_s \) is the velocity of particle in the region \( z < 0 \) and \( \overline{\psi_{tr}} \) is the complex conjugate of \( \psi_{tr} \).
\[
T = \frac{A \left( \frac{1}{2} \right) \left( \frac{1}{2} \frac{k_2}{k_1} \right) e^{ik_2z} - \frac{A}{2i} \left( \frac{1}{2} \frac{k_2}{k_1} \right) e^{-ik_2z} \frac{A \left( \frac{1}{2} \right) \left( \frac{1}{2} \frac{k_2}{k_1} \right) e^{-ik_2z} - \frac{A}{2i} \left( \frac{1}{2} \frac{k_2}{k_1} \right) e^{ik_2z}} {2 \left( \frac{1}{2} \frac{1}{2} \frac{k_2}{k_1} \right) e^{ik_1z} - 2 \frac{1}{2} \frac{1}{2} \frac{k_2}{k_1} e^{-ik_1z}} \]
Or \( T = \frac{\left( 1 - \frac{k_2}{k_1} \right) \left( 1 - \frac{k_2}{k_1} \right) e^{-ik_1z}} {1 + \frac{k_2}{k_1} \left( 1 - \frac{k_2}{k_1} \right) e^{ik_1z}} \)
On simplifying, we get
\[
T = \frac{\left| k_1 - k_2 \right|^2} {\left| k_1 + k_2 \right|^2} \ldots \ldots \ldots \ldots \ldots (10) \]
Using the relation \[^{[10]}\] between the momentum (or velocity) and the propagation constant, we can write
\[
p_2 = \frac{mv_2}{k_2}, \quad p_1 = \frac{mv_1}{k_1} \]
Or \( \frac{v_2}{v_1} = \frac{k_1}{k_2} \)
Hence we can re write equation (10) as
\[
T = \frac{4k_2}{k_1 \left( 1 + \left| k_2 \right|^2 \right)} \ldots \ldots \ldots \ldots \ldots (11) \]
The equation (11) provides the expression for the quantum mechanical transmission coefficient at the vertical step potential with \( E > V_0 \). On adding equations, (9) and (11), we can find that the sum of \( R \) and \( T \) is one which shows that the sum of the number of electrons reflected at the potential discontinuity and the number of electrons transmitted through the potential discontinuity is equal to number of electrons incident on the potential discontinuity with \( E > V_0 \).

**Case II: \( E < V_0 \)**

In this case, in the region, \( z < 0 \), the solution remains the same. In the region, \( z > 0 \), \( V(z) = V_0 \), \( E < V_0 \), therefore, the time – independent Schrodinger equation in one dimensions is written as
\[
D_2^2 \Psi_R(z) + k_3^2 \Psi_R(z) = 0 \ldots (12), \]
where
\[
k_3 = \sqrt{\frac{2m(V_0 - E)}{h^2}}, \quad k = \sqrt{\frac{2m(V_0 - E)}{h^2}} \]
Taking Laplace transforms of equation (12), we get
\[
r^2 \psi_R(r) - r \psi_R(0) - D_2 \psi_R(0) + k_3^2 \psi_R(r) = 0, \]
where \( \psi_R(r) \) is the Laplace transform of \( \psi_R(z) \).
Or
\[
r^2 \psi_R(r) - r A - B + k_3^2 \psi_R(r) = 0. \]
Rearranging the equation, we get
\[
\psi_R(r) = \frac{r A + B}{r^2 + k_3^2} \]
Or \( \overline{\psi_R(r)} = \frac{r A - B}{r^2 + k_3^2} \)
Taking inverse Laplace transform, we get
\[
\Psi_R(z) = A \cos k_3 z + \frac{B}{k_3} \sin k_3 z \ldots \ldots \ldots \ldots \ldots (13) \]
Writing \( \cos k_3 z \) and \( \sin k_3 z \) in terms of exponentials, we get
\[
\Psi_R(z) = A \frac{e^{ik_3 z} + e^{-ik_3 z}}{2} + B \frac{e^{ik_3 z} - e^{-ik_3 z}}{2i} \]
Or
\[
\Psi_R(z) = \left( \frac{A}{2} + \frac{B}{2i k_3} \right) e^{ik_3 z} + \left( \frac{A}{2} - \frac{B}{2i k_3} \right) e^{-ik_3 z} \]
Or
\[
\Psi_R(z) = \left( \frac{A}{2} - \frac{B}{2k_3} \right) e^{ik_3 z} + \left( \frac{A}{2} + \frac{B}{2k_3} \right) e^{-ik_3 z} \ldots \ldots \ldots \ldots \ldots (14) \]
In equation (14), \( \left( \frac{A}{2} - \frac{B}{2k_3} \right) e^{ik_3 z} \) represents the transmitted wave in the region \( z > 0 \) i.e.
\[
\psi_{tr}(z) = \left( \frac{A}{2} - \frac{B}{2k_3} \right) e^{ik_3 z} \ldots \ldots \ldots \ldots \ldots (15) \]
Quantum Mechanical Reflection and Transmission Coefficients for a Particle through a One-Dimensional Vertical Step Potential

Since \( \left( \frac{A}{2} + i \frac{B}{2k_3} \right) \) represents the coefficient of a beam incident from right on the potential step, which is not physical, therefore, \( \frac{A}{2} + i \frac{B}{2k_3} = 0 \)

Or \( \frac{A}{2} + i \frac{B}{2k_3} = 0 \)

Or \( B = ik_3A \ldots (16) \)

Using equations (16), we can write

\[
\psi_{in}(z) = \frac{A}{2} \left( 1 + \frac{k_3}{k_1} \right) e^{ik_1z},
\]

\[
\psi_{re}(z) = \frac{A}{2} \left( 1 - \frac{k_3}{k_1} \right) e^{-ik_1z} \quad \text{and}
\]

\[
\psi_{tr}(z) = A e^{ik_3z}.
\]

The quantum mechanical reflection coefficient \( R \) is given by [9-9]

\[
R = \frac{\text{Reflected probability flux}}{\text{Incident probability flux}}
\]

Or \( R = \frac{\psi_{re} \bar{\psi}_{re}}{\psi_{in} \bar{\psi}_{in}} \), where \( \psi_i \) is the velocity of particle in the region \( z < 0 \) and \( \psi_{re}^* \) is the complex conjugate of \( \psi_{re} \) and \( \psi_{in}^* \) is the complex conjugate of \( \psi_{in} \).

Or \( R = \frac{\frac{1}{2} \left( 1 + \frac{k_3}{k_1} \right) e^{-ik_1z}}{\frac{1}{2} \left( 1 - \frac{k_3}{k_1} \right) e^{ik_1z}} \), where \( A^* \) and \( k_3^* \) are the complex conjugates of \( A \) and \( k_3 \).

Or \( R = \frac{k_3}{k_1} \left( 1 + \frac{k_3}{k_1} \right) \cdot \left( 1 + \frac{k_3}{k_1} \right)^{-1} \)

Or \( R = 1 \ldots (17) \)

The equation (17) provides the expression for the quantum mechanical reflection coefficient at the vertical step potential with \( E < V_0 \).

The quantum mechanical transmission coefficient \( T \) is given by [9-9]

\[
T = \frac{\text{Transmitted probability flux}}{\text{Incident probability flux}}
\]

Or \( T = \frac{\psi_{tr} \bar{\psi}_{tr}}{\psi_{in} \bar{\psi}_{in}} \), where \( \psi_2 \) is the velocity of particle in the region \( z < 0 \) and \( \psi_{tr}^* \) is the complex conjugate of \( \psi_{tr} \).

\[
T = \frac{v_2 A e^{-ik_3z} \bar{A} e^{ik_3z}}{v_1 A e^{-ik_3z} \bar{A} e^{ik_3z}}
\]

\[
T = \frac{v_2 A e^{-ik_3z} \bar{A} e^{ik_3z}}{v_1 A e^{-ik_3z} \bar{A} e^{ik_3z}}
\]

Or \( T = \frac{4v_2}{v_1} \left( 1 + \frac{k_3}{k_1} \right)^{1/2} \left( \frac{k_1}{k_1} \right)^{-1/2} e^{-ik_1z} \)

Using the relation [10] between the velocity and the propagation constant, we can write

\[
\frac{v_2}{v_1} = \frac{k_3}{k_1}
\]

Or \( \frac{v_2}{v_1} = \frac{k_3}{k_1} \)

Hence we can rewrite equation (18) as

\[
T = \frac{4k_3}{k_1} \left( 1 + \frac{k_3}{k_1} \right)^{-1}, \quad \text{which is complex whose real part is given by}
\]

\[
\text{Real}(T) = 0 \ldots (19)
\]

The equation (9) provides the expression for the quantum mechanical reflection coefficient at the vertical step potential with \( E < V_0 \).

Again, on adding equations (17) and (19), we can find \( R + T = 1 \) which shows that the sum of the number of electrons reflected at the potential discontinuity and the number of electrons transmitted through the potential discontinuity is equal to number of electrons incident on the potential discontinuity with \( E < V_0 \).

IV. RESULT AND DISCUSSION

In this paper, we have discussed the quantum mechanical Reflection and Transmission coefficients for a particle through a one-dimensional vertical step potential through the Laplace transformation method. It is revealed that the quantum mechanical Reflection and Transmission coefficients for a particle through a one-dimensional vertical step potential can be obtained easily through the application of Laplace transformation. This method results out a new technique in discussing the quantum mechanical Reflection and Transmission coefficients for a particle through a one-dimensional vertical step potential, and it served well to illustrate the application of the Laplace transformation tool.

V. CONCLUSION

It is concluded that sum of \( R \) and \( T \) (whether \( E > V_0 \) or \( E < V_0 \)) is one which shows that the sum of the number of electrons reflected at the potential discontinuity (i.e. the flux of reflected particles) and the number of electrons transmitted through the potential discontinuity (i.e. the flux of transmitted particles) is always equal to number of electrons incident on the potential discontinuity (i.e. the flux of incident particles) which means that no particle is lost at the potential discontinuity.

REFERENCES

1. Quantum mechanics by B.N. Srivastava, R.M. Richaria. 16th edition, 2017.
2. Principles of quantum mechanics by P.A.M. Dirac. Edition: Reprint, 2016.
3. Principles of quantum mechanics by R. Shankar. 2nd edition
4. Quantum mechanics by Satya Parkash. Edition: Reprint, 2016.
5. Advanced Engineering Mathematics by Erwin Kreysig 10th edition, 2014.
6. Advanced engineering mathematics by H.K. Dass. Edition: Reprint, 2014.
7. Gupta. R., Gupta. R, Verma. D., ‘Eigen Energy Values and Eigen Functions of a Particle in an Infinite Square Well Potential by Laplace Transform’, International Journal of Innovative Technology and Exploring Engineering (IJITEE), Volume-8 Issue-3, January.
8. Quantum Mechanics and Path Integrals by Richard P. Feynman, Albert R. Hibbs, Daniel F. Styer.
9. Introduction to Quantum mechanics by David j. Griffiths. Edition 2nd, 2017
10. A text book of Engineering Physics by M.N. Avadhanulu. Revised edition 2014
11. Higher Engineering Mathematics by Dr.B.S.Grewal 43rd edition 2015.
12. Introduction to Laplace transforms and Fourier series by Dyke and Phil.
13. Essential mathematical methods by K. F. Riley and M.P. Hobson.
AUTHORS PROFILE

Mr. Rohit Gupta is currently Lecturer, Department of Physics (Applied Sciences), Yogananda College of Engineering and Technology, Gusha Brahmana (Patoli), Akhnoor Road, Jammu (J&K, India). He has done M.Sc. Physics from University of Jammu (J&K) in the year 2012. He has been teaching UG Classes for well over six and half years. He has to his credit Sixteen Research Papers. He has attended three Workshop/Conferences/FDP during his six and half years’ experience of teaching.

Dr. Tarun Singhal is currently Assistant Professor, Department of Applied Sciences, CGC College of Engineering, Landran Mohali, Chandigarh, India. He has done Ph.D. in Electronics and Communication Engineering from Maharishi Dayanand University Rohtak Haryana, India. He has been teaching UG and PG Classes for well over thirteen years. He has to his credit Eleven Research Papers. He has attended Twenty five Workshop/Conferences/FDP during his Thirteen years’ experience of teaching.

Dr. Dinesh Verma is currently Associate Professor, Department of Mathematics, Yogananda College of Engineering and Technology, Gurha Brahmana (Patoli), Akhnoor Road, Jammu (J&K, India). He has done Ph.D. at M.J.P. Rohilkhand University, Bareilly (U.P.) in 2004. He has been teaching UG Classes for well over two decades. He has to his credit four books for Engineering and Graduation level which are used by students of various universities. He has to his credit Twenty Two Research Papers. He has attended Seventeen Workshop/Conferences/FDP during his Twenty years’ experience of teaching. Also, he has a membership with ISTE (Indian Society for Technical Education) and ISCA (Indian Science Congress Association).