You Can’t Always Get What You Want: The Impact of Prior Assumptions on Interpreting GW190412

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Abstract

GW190412 is the first observation of a black hole binary with definitively unequal masses. GW190412’s mass asymmetry, along with the measured positive effective inspiral spin, allowed for inference of a component black hole spin: the primary black hole in the system was found to have a dimensionless spin magnitude between 0.17 and 0.59 (90% credible range). We investigate how the choice of priors for the spin magnitudes and tilts of the component black holes affect the robustness of parameter estimates for GW190412, and report Bayes factors across a suite of prior assumptions. Depending on the waveform family used to describe the signal, we find either marginal to moderate (2:1–6:1) or strong (≥20:1) support for the primary black hole being spinning compared to cases where only the secondary is allowed to have spin. We show how these choices influence parameter estimates, and find the asymmetric masses and positive effective inspiral spin of GW190412 to be qualitatively, but not quantitatively, robust to prior assumptions. Our results highlight the importance of both considering astrophysically motivated or population-based priors in interpreting observations and considering their relative support from the data.

Unified Astronomy Thesaurus concepts: Gravitational waves (678); LIGO (920); Gravitational wave sources (677); Binary stars (154); Bayesian statistics (1900); Black holes (162); Stellar evolution (1599)

1. Introduction

GW190412 (Abbott et al. 2020) was the first reported observation of a binary black hole (BBH) from the third observing run (O3) of the Advanced LIGO (Aasi et al. 2015) and Advanced Virgo (Acernese et al. 2015) detector network. GW190412’s source is the first system to have definitively unequal masses (see Abbott et al. 2019a), with the primary (BH) being ~30 $M_\odot$ and the secondary BH being ~8 $M_\odot$. In addition to unveiling emission from higher-order multipoles (HMs), this asymmetry allowed for enhanced constraints on the intrinsic and extrinsic parameters of the BBH system.

The spins of compact binary components are difficult to measure from gravitational-wave (GW) signals (Poisson & Will 1995; Vitale et al. 2014; Abbott et al. 2016a; Pürrer et al. 2016). Typically, spin constraints are presented in terms of mass-weighted combinations of the two component spins: the effective inspiral spin

$$\chi_{\text{eff}} = \frac{m_1 \chi_1 \cos \theta_1 + m_2 \chi_2 \cos \theta_2}{m_1 + m_2},$$

where $m_1 \geq m_2$ are the component masses, $\chi_i$ are the dimensionless spin magnitudes, and $\theta_i$ are the angles between the spins and the Newtonian orbital angular momentum, $\mathbf{L}$, encodes information about the spin components aligned with the orbital angular momentum (Damour 2001; Racine 2008; Santamaria et al. 2010; Ajith et al. 2011), whereas in-plane spins are characterized by the effective precession spin

$$\chi_p = \max \left\{ \chi_1 \sin \theta_1, \frac{q (4q + 3)}{4 + 3q} \chi_2 \sin \theta_2 \right\}.$$

The LIGO Scientific & Virgo Collaboration (LVC) reported an effective spin for GW190412 of $\chi_{\text{eff}} = 0.25_{-0.11}^{+0.08}$ (median and 90% credible interval; Abbott et al. 2020). Since $\chi_{\text{eff}}$ is positive and constrained away from zero, at least one of the BHs in the GW190412 system had a spin direction in the same hemisphere as $\mathbf{L}$ during the GW inspiral. GW190412 also exhibited marginal hints of orbital precession, which is consistent with at least one of the BH spins being nonzero.

A BBH with $\chi_{\text{eff}} > 0$ has been observed before in GW151226 (Abbott et al. 2016b; Miller et al. 2020), and potentially in GW170729 (Abbott et al. 2019a; Chatziioannou et al. 2019). However, the larger mass of the primary BH in GW190412 relative to the secondary BH allowed for the spin of the primary to be inferred as $\chi_1 = 0.43^{+0.16}_{-0.26}$. This is because when $m_1 \gg m_2$ the primary spin is much more important in determining the dynamics of the system (as illustrated by the mass weighting in $\chi_{\text{eff}}$ and $\chi_p$), and we are less sensitive to the value of the secondary spin. GW190412 therefore is the first high-significance detection of a compact binary system with an observable component spin.
GW190412’s primary spin may be difficult to reconcile with theoretical modeling of massive binary stars in isolation. Detailed modeling of core–envelope interaction in massive stars finds angular momentum transport to be highly efficient, driving stellar cores to extremely slow rotation prior to their collapse into a compact object (Kushnir et al. 2016; Zaldarriaga et al. 2018; Fuller et al. 2019; Fuller & Ma 2019; Belczynski et al. 2020). This theoretical underpinning is corroborated by current GW catalogs, which contain systems that are mostly consistent with \( \chi_{\text{eff}} \approx 0 \) (Abbott et al. 2019b; Miller et al. 2020; Nitz et al. 2020; Vennumadhav et al. 2020). Though the birth spins of some BHs in high-mass X-ray binaries have been interpreted as near extremal (\( \chi \approx 1 \); see Miller & Miller 2015 and references therein), it is unclear whether these systems will evolve to be BBHs that merge within a Hubble time (e.g., Belczynski et al. 2012; Qin et al. 2019). Following this reasoning, multiple groups have proposed that the high spin of the primary BH in GW190412 is the result of an alternative formation scenario to canonical isolated binary evolution, such as dynamical assembly in young star clusters (Di Carlo et al. 2020), hierarchical mergers in massive stellar clusters (Gerosa et al. 2020; Kimball et al. 2020; Rodriguez et al. 2020), active galactic nucleus (AGN) disks (Tagawa et al. 2020), Population III stars (Kinugawa et al. 2020), and mergers induced from the secular evolution in hierarchical systems (Hamers & Safarzadeh 2020).

On the other hand, the second-born BH in BBH merger progenitors can be significantly spun up through tidal locking of the stellar core with the first-born BH (Qin et al. 2018; Bavera et al. 2020). If GW190412 could instead be explained by a highly spinning secondary BH, the standard isolated formation scenario with a low-spinning primary could again be viable. To this end, Mandel & Fragos (2020) provide a reinterpretation of the LVC analysis (Abbott et al. 2020) using a prior motivated by theoretical predictions of BBH progenitors formed in isolation. Assuming a prior with a zero-spin primary BH and a secondary BH whose spin projection is aligned with the orbital angular momentum, Mandel & Fragos (2020) reweight the public posterior samples of GW190412 (LIGO Scientific Collaboration & Virgo Collaboration 2020), effectively interpreting the measured value of \( \chi_{\text{eff}} \) as originating from the secondary’s spin rather than the primary’s. To compensate for the nonzero effective spin of GW190412, the reweighted posteriors from this analysis point to a highly spinning secondary BH with \( \chi_2 \gtrsim 0.64 \) (Mandel & Fragos 2020).

Though predictions for the formation rate of these systems are highly sensitive to the uncertain prescription for natal BH spins, recent work has found that for systems with asymmetric masses such as GW190412, the highly spinning secondary BH interpretation is more probable from an isolated evolution standpoint than a moderately spinning primary (e.g., Olejak et al. 2020). This is consistent with the current catalog of GWs, since individual spins are poorly constrained in all previously observed BBHs (Abbott et al. 2019b). However, even this formation mechanism struggles to accommodate GW190412, as systems where the secondary BH has been significantly spun up due to tidal interactions have short merger timescales and a merger rate in the local universe that is at least an order of magnitude lower than what is estimated for GW190412-like systems (Safarzadeh & Hotokezaka 2020).

Nonetheless, while various assumptions may be made to represent the prior belief for parameters given an astrophysical model, it is critical to determine whether a given model is supported by the data. The amount by which the data supports a specific model (in this work, a prior) is encoded in the Bayesian evidence. While varying prior assumptions will yield differing parameter estimates, the ratio of evidences between models—the Bayes factor \( B \)—indicates whether any one prior assumption is favored or disfavored by the data compared to another. This is particularly important to verify for the case of strong priors, since they might drive the posteriors to potentially arbitrary values at the expense of the evidence: if you torture the data long enough, it will confess to anything (Coase 1982).

For example, in the analysis of GW151226 (Abbott et al. 2016b), Vitale et al. (2017a) showed how if one uses a prior that enforces small (~0.1) spin magnitudes, the evidence decreases by a factor of 50 compared to a uniform prior, while the posteriors still look reasonable. It is only by comparing evidences between models, i.e., calculating Bayes factors, that one can assess which model is better described by the data.

In this Letter, we explore various prior assumptions for the interpretation of GW190412 and calculate Bayes factors between these model assumptions. The priors we choose are motivated by various astrophysical models presented in the literature, with a particular focus on the spin of the second-born BH, and the astrophysically relevant question of whether the primary is spinning.

In Section 2 we explain the various prior assumptions we choose when analyzing the data, and their astrophysical motivation. We present Bayes factors across these prior assumptions in Section 3, and examine the impact of differing prior assumptions on the parameter estimation for GW190412 in Section 4. In Section 5 we discuss the results of our analysis and their impact on the interpretation of GW190412, and comment on astrophysical implications.

### 2. Data Analysis and Prior Assumptions

To investigate the impact of prior assumptions on the inferred parameters of GW190412 and the Bayes factors between these assumptions, we perform parameter estimation using a suite of prior assumptions motivated by various astrophysical predictions. We use the publicly available data for GW190412 (LIGO Scientific Collaboration & Virgo Collaboration 2020) and follow the parameter-estimation procedure used in Abbott et al. (2020). Our results are produced using a highly parallelized version of BILBY (Ashton et al. 2019; Smith et al. 2019; Romero-Shaw et al. 2020), which computes posterior probability distributions for the properties of the source as well as model evidence.

We use both the Phenom and EOB families of waveform approximants in our analysis.7 We use IMRPhenomPv3HM (Khan et al. 2019, 2020) and SEOBNRv4PHM (Pan et al. 2014; Babak et al. 2017; Ossokine et al. 2020), both of which include the effects of spin precession and HM modes. Inclusion of HMs in waveform approximants is crucial for the parameter estimation of GW190412, as this more complete physical picture of the GW signal is necessary to accurately constrain the mass ratio \( q = m_2/m_1 \) and spins (Van Den Broeck & Sengupta 2007; Graff et al. 2015; Calderón Bustillo et al. 2016; Varma & Ajith 2017; Abbott et al. 2020). Systematic

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7 Different waveform models are referred to as “approximants” throughout to differentiate between waveform approximant models and the models representing different prior configurations.
differences are expected between analyses using Phenom and EOB approximants, as evident in Abbott et al. (2020). Though we use the Phenom approximant for all seven prior configurations described below, due to the computational cost of the EOB approximant we only run this with two exemplary prior configurations.

The priors we consider are:

(A) Uniform in spin magnitude for both components, isotropic and unconstrained in spin tilts. This uninformative prior is used in Abbott et al. (2020); it does not make strong assumptions about spin orientations or magnitudes, and its broad support enables reweighting by different priors (e.g., Mandel et al. 2019; Thrane & Talbot 2019).

(B) Uniform in spin magnitude and isotropic in spin tilt for the primary BH, with a non-spinning secondary. A spinning primary and a non-spinning secondary may be expected if BHs are born with small spins, but the larger BH is the result of a previous BH merger and has gone on to form a new binary in a dense stellar environment such as a globular or nuclear cluster (Fishbach et al. 2017; Gerosa & Berri 2017; Rodriguez et al. 2019; Gerosa et al. 2020; Kimball et al. 2020). In this scenario, we would typically expect the primary spin magnitude to be $\chi_1 \sim 0.67$.

(C) Non-spinning primary BH with unconstrained spin for the secondary BH. This is representative of an isolated formation scenario, with a secondary that can be spun up through tidal interactions (Qin et al. 2019; Bavera et al. 2020). The unconstrained spin tilt, however, allows for significantly misaligned spins, which are difficult to attain for BBHs in the standard isolated evolution scenario (e.g., Kalogera 2000; Fryer et al. 2012; Rodriguez et al. 2016).

(D) Same as Model C, but with spin tilts constrained to be in the same hemisphere as the orbital angular momentum: $\chi_2 \cdot \hat{\mathbf{L}} \geq 0$ ($0^\circ \leq \theta_2 \leq 90^\circ$). This is similar to the prior assumption used in Mandel & Fragos (2020).

(E) Same as Model C, but with spin tilts for the secondary constrained to $0^\circ \geq \theta_2 \geq 10^\circ$. This model has been used to represent near-aligned spins (e.g., Vitale et al. 2017b), as predicted from the coevolution of isolated binaries and weak BH natal kicks at birth.

(F) Same as Model C, but with spin tilts for the secondary perfectly aligned with the orbital angular momentum ($\theta_2 = 0^\circ$).

(G) Non-spinning primary and secondary. This is an extreme assumption that we expect will struggle to match the data due to the positive measured $\chi_{\text{eff}}$ and marginal precessional information.

These configurations are summarized on the left side of Table 1. For all other parameters, we use priors analogous to those used by the LVC in the analysis of GW190412 (Abbott et al. 2020).

### Table 1

| Model | Prior Assumption | Evidence |
|-------|-----------------|----------|
| A     | $\chi_1$ \ U[0, 0.99], $\theta_1$ \ $\chi_2$ \ $\theta_2$ | $\log_{10} C_{\text{max}}$, $B$, $\log_{10}(B)$, $\delta_{\log_{10}(B)}$ |
| A-EOB | $\chi_1$ \ $\chi_2$ \ $\theta_1$ \ $\theta_2$ | 77.0, 1.0, 0.00, 0.08 |
| B     | $\chi_1$ \ $\chi_2$, $\theta_1$ \ $\theta_2$ | 77.1, 1.0, 0.00, 0.10 |
| C     | $\delta(0)$ \ $\chi_1$ \ $\chi_2$ \ $\theta_1$ \ $\theta_2$ | 76.5, 6.2 \ $10^{-1}$, –0.20, 0.09 |
| D     | $\delta(0)$ \ $\chi_1$ \ $\chi_2$, $\theta_1$ \ $\theta_2$ | 75.6, 1.5 \ $10^{-1}$, –0.80, 0.09 |
| D-EOB | $\delta(0)$ \ $\chi_1$ \ $\chi_2$ \ $\theta_1$ \ $\theta_2$ | 75.5, 3.6 \ $10^{-1}$, –0.44, 0.09 |
| E     | $\delta(0)$ \ $\chi_1$ \ $\chi_2$, $\theta_1$ \ $\theta_2$ | 74.4, 4.9 \ $10^{-2}$, –1.30, 0.10 |
| F     | $\delta(0)$ \ $\chi_1$ \ $\chi_2$, $\theta_1$ \ $\theta_2$ | 75.3, 5.3 \ $10^{-1}$, –0.27, 0.09 |
| G     | $\delta(0)$ \ $\chi_1$ \ $\chi_2$, $\theta_1$ \ $\theta_2$ | 75.2, 5.0 \ $10^{-1}$, –0.30, 0.09 |
|       |                 | 71.5, 2.8 \ $10^{-3}$, –2.54, 0.08 |

Note. We use short-hand for the distributions we consider: $\delta$ uses a fixed value, $U$ denotes a uniform distribution, and ISO is an isotropic distribution (uniform in $\cos(\theta)$). Angular assumptions are omitted when the spin magnitudes for that component are forced to zero. The right shows maximum value of the log-likelihood ($\log_{10} C_{\text{max}}$) and Bayes factors ($B$) for each model. Each Bayes factor is calculated relative to the uninformative LVC prior for the respective waveform approximant (Model A and Model A-EOB for Phenom and EOB, respectively). For reference, $B < 1$ ($\log_{10} B < 0$) means that data prefer the reference model, $B \geq 3:1$ ($\log_{10} B \geq 0.5$) indicates moderate evidence for the new hypothesis, and $B \geq 10:1$ ($\log_{10} B \geq 1.0$) indicates strong evidence for the new hypothesis. The rightmost column gives the estimated 1σ uncertainty in $\log_{10} B$.
describing our
is the likelihood
could also be used to compare waveform approximants; for compared to the standard LVC prior that only includes the effects of the spin components aligned with the orbital between using a waveform approximant that includes spin precession and one
Analogously, when spin precession is not measurable, such that the evidences for the two models.
The model evidence, or marginalized likelihood, is
where the integral is over the parameters \( \theta \), describing our source (masses, spins, etc.), \( L(\theta) = p(d|\theta) \) is the likelihood of the parameters (Cutler & Flanagan 1994), and \( p(\theta|M_i) \) is our prior probability density on the parameters within model \( M_i \), as described in Section 2. Thus, the Bayes factor is given by
When considering models with more parameters, or with parameters allowed to vary on a larger domain, we expect that we may be able to fit the data better, giving higher likelihoods. In calculating evidences, this is counterbalanced by the increased prior volume: as we spread the total prior probability (which must integrate to 1) over a larger volume where the likelihood can have potentially negligible support, its density around the maximum likelihood region may decrease, resulting in a lower evidence. This Occam factor allows the Bayes factor to be used to determine if more complicated models are needed to explain data (MacKay 2003, Chapter 28). When considering spins measured with GW observations, we are typically only sensitive to particular mass-weighted combinations of the 6 spin degrees of freedom (Poisson & Will 1995; Chatziioannou et al. 2014; Pürrer et al. 2016; Vitale et al. 2017a). Therefore, it may be possible to fit the data well by assuming only a single component is spinning, and we would not anticipate a strong preference in favor of a more complicated model including two spinning bodies. In cases where there is a large asymmetry in masses the secondary spin may become irrelevant, and the properties of the signal may be completely determined by the primary spin. When the secondary spin has negligible impact on the likelihood, we expect there will be no preference between models with and without a secondary spin as it is unconstrained and its introduction incurs no Occam factor penalty.
In Table 1 we show Bayes factors for each prior configuration compared to the standard LVC prior (Model A). Bayes factors could also be used to compare waveform approximants; for example, the Bayes factor between the Phenom and EOB approximants using the LVC prior (Models A and A-EOB) is 0.97:1, indicating no preference for one of these approximants over the other. Since we focus on the impact of differing prior

4. Parameter Estimation
Prior assumptions can have a strong effect on the measurement of intrinsic and extrinsic parameters inherent to a BBH coalescence. Here, we investigate the robustness of parameter estimates for GW190412 across our various prior assumptions, with a particular focus on spin parameters.

4.1. Mass Ratio
GW190412 is the first BBH with definitely unequal masses, with a reported mass ratio at the 90% credible level of 0.25 \( \leq q \leq 0.45 \) using the Phenom approximant and 0.21 \( \leq q \leq 0.31 \) using the EOB approximant (Abbott et al. 2020). In Figure 1 we show the posterior distributions for \( q \) across our different priors and waveform approximants. Aside from the (strongly disfavored) Model G, which does not allow
for spins in either BH, we find the mass ratio to be constrained to $q \leq 0.57$ at the 99% credible level.

There is a noticeable difference in the posterior distribution for $q$ when using priors where the primary is spinning compared to those where only the secondary is spinning. We find that the posterior for $q$ pushes to larger values when $\chi_1 = 0$, with a median of 0.39 (0.39) in Model D compared to 0.32 (0.26) in Model A when using the Phenom (EOB) approximant. This change in $q$ results in a more massive secondary that can more easily account for the observed effective spins.

### 4.3. Component Spins

In Figure 2 we show marginalized posterior distributions for the two component spins, $\chi_1$ and $\chi_2$. In the prior configurations where $\chi_1$ is nonzero, we recover similar posterior distributions across the Phenom results, though when $\chi_2$ is forced to zero the distribution shifts to slightly higher values with a median $\chi_2$ that is 0.03 larger than in the LVC prior case. This is because the primary BH must now account for all spin effects in the data without a contribution from the secondary. The $\chi_1$ posteriors are also consistent with the Bayes factors reported in Table 1 in favor of prior configurations where the primary BH is spinning: for Models A and C (which are nested since Model C can be obtained by fixing $\chi_1 = 0$ in Model A), the Bayes factor can also be calculated by comparing the prior to the posterior at $\chi_1 = 0$ (Chatziioannou et al. 2014). As evident in the top panel of Figure 2, the prior at $\chi_1 = 0$ is larger than the posterior for both waveform approximants, pointing to a Bayes factor in favor of a spinning primary. Estimating Bayes factors from the posterior and prior densities is subject to considerable sampling error when considering the tails of the distributions.
where there are few samples. Nonetheless, for the Phenom approximant, we find the prior probability density at $\chi_2 = 0$ in Model A to be a factor of $\approx 5.5$ larger than the posterior probability density at $\chi_2 = 0$, which is consistent with the Bayes factor between these two prior configurations ($\approx 6.4:1$).

We see larger variation in $\chi_2$ across the prior configurations. The standard LVC prior recovers a broad, uninformative distribution in $\chi_2$. However, when $\chi_2$ is forced to zero spin, $\chi_2$ is constrained away from zero in all cases; in these prior configurations, we find $\chi_2$ to be consistent with maximally spinning and have $\chi_2 \gtrsim 0.50$ at the 90% credible level (Mandel & Fragos 2020). The EOB results push to slightly higher secondary spins than the Phenom results with non-spinning primary configurations, with $\chi_2 \gtrsim 0.62$ at the 90% credible level for Model D-EOB. In all cases where spin misalignment is allowed, we find a preference for some degree of misalignment in the spins; at the 90% credible level, we find $\theta_2 > 18.5^\circ$ in Model C, $\theta_2 > 18.1^\circ$ ($\theta_2 > 16.3^\circ$) in Model D (Model D-EOB), and $\theta_2 > 3.2^\circ$ in Model E. Thus, we find that precession (albeit possibly immeasurable) is permitted in all prior configurations that allow for spin tilts.

### 5. Discussion and Conclusions

GW190412 is an astrophysically compelling event that resides in a previously unobserved region of BBH parameter space. The effective inspiral spin of the system indicates that at least one of the component BHs is spinning. This work investigates whether a spinning primary BH or a spinning secondary BH is better supported by the data, and how these hypotheses affect the inferred parameters of GW190412.

Our main results are summarized in Table 1. The broad LVC prior (Model A) with both BHs spinning is preferred over the other prior configurations, despite the larger prior volume. The degree of preference depends on the waveform approximant used, as the effect of waveform systematics are nonnegligible for this event (Abbott et al. 2020). We recover marginal support in favor of Model A compared to the prior configuration where only the primary is spinning (Model B). When using the Phenom approximant we find marginal to moderate evidence in favor of Model A compared to prior configurations where only the secondary is allowed to spin (Models C–F), whereas with the EOB approximant we find strong evidence in support of the LVC prior configuration compared to priors with a non-spinning primary. The data strongly support Model A over the hypothesis where neither BH is spinning (Model G) for both waveform approximants.

The Phenom approximant gives broader parameter constraints than the EOB approximant in both Abbott et al. (2020) and this work. In Figure 1, we see that the non-spinning primary prior configurations move the posterior distributions for $q$ ($\chi_{\text{eff}}$) to higher (lower) values to better allow the secondary to account for the spin information in the signal. This comes at the cost of matching the data, as the maximum likelihood values are $\lesssim 1.5$ dex lower for prior configurations where only the secondary is spinning compared to Model A. Whereas the Phenom approximant measures $q$ ($\chi_{\text{eff}}$) to be $\lesssim 0.45$ ($0.29$) and the 90% credible level, the EOB approximant recovers $\lesssim 0.31$ ($0.34$). The lower mass ratio and higher effective spin from the EOB analysis makes it more difficult for the data to accommodate a non-spinning primary, with maximum likelihood values that are $\approx 2.7$ dex lower for the prior configuration where only the secondary is spinning. Despite the larger prior volume, we find a spinning primary hypothesis to be favored over a non-spinning primary hypothesis by a Bayes factor of $\gtrsim 20:1$.

The prospect of a non-spinning primary BH was explored in Mandel & Fragos (2020). Mandel & Fragos (2020) reweighted the publicly released $\chi_{\text{eff}}$ posterior samples from the EOB analysis (Abbott et al. 2020) in order to apply a prior that assumes a non-spinning primary and a secondary that has a spin aligned with the orbital angular momentum. This approach assumes that there is a single measurable spin degree of freedom from GW190412 that is identified with $\chi_{\text{eff}}$, and that there is no information about spin precession. We instead reanalyze the data under the desired prior, thus imposing no such restrictions about how spins are measured. Our analysis results in similar constraints on the secondary spin (Model D with EOB) as Mandel & Fragos (2020), but a different estimate of the mass ratio; we find $0.34 \leq q \leq 0.47$ at the 90% level, compared to $0.27 \leq q \leq 0.36$. This difference could be attributed to the assumptions of Mandel & Fragos (2020) about spin measurability. For example, we find that the data contain small (but nonnegligible) information about spin precession. Additionally, the leading-order spin term in the GW phase is not identical to $\chi_{\text{eff}}$ (Poisson & Will 1995), with the difference between the two being more prominent for unequal mass systems such as GW190412. This suggests that the relation between $\chi_2$ and $q$ cannot be fully explored when considering only $\chi_{\text{eff}}$. Both Mandel & Fragos (2020) and this study conclude that the assumption of a non-spinning primary requires a highly spinning secondary, although we find that the corresponding Bayes factors disfavor this scenario.

Regardless of our prior assumptions, we find the positive effective spin and unequal masses of GW190412 to be robust conclusions. However, we do see a shift in the posterior distributions across our prior assumptions. With only the secondary BH spinning, we recover higher values for $q$ and lower values for $\chi_{\text{eff}}$ and $\chi_q$. The component spins are affected more dramatically; forcing a non-spinning primary causes the secondary’s spin magnitude posterior to significantly increase and shift against the physical boundary at $\chi_2 = 1$.

The sensitivity of parameter-estimation results to the choice of prior highlights the importance of choosing an appropriate prior when interpreting observations. One will never find a spinning primary BH if spins are always restricted to be zero; astrophysical models are uncertain and need to be constrained by observations. To this end, we can construct prior distributions using a population of observations. Performing hierarchical inference enables inference of both individual event’s properties and those of the population (Mandel 2010; Abbott et al. 2019b; Gallaudage et al. 2019), in effect using the set of observations to construct an empirical prior. These inferences may use a branching fraction to consider models from different formation channels (Stevenson et al. 2017; Talbot & Thrane 2017; Vitale et al. 2017b; Zevin et al. 2017) or use a phenomenological model to describe the underlying population (Roulet & Zaldarriaga 2018; Fishbach et al. 2018; Wysocki et al. 2019; Fishbach et al. 2020); they may even encode prior odds for different channels (Kimball et al. 2020). Using wide, uninformative priors, as done by the LVC, enables parameter-estimation results to be reweighted by different priors, as required for a hierarchical population analysis (Mandel et al. 2019; Thrane & Talbot 2019).
Both the moderately spinning primary and highly spinning secondary interpretations for GW190412 provide unprecedented constraints on astrophysical formation scenarios. If GW190412 is the product of isolated binary evolution, our results indicate that the paradigm of negligible natal spin for the first-born BH in BBH merger progenitors may need to be revised (Kushnir et al. 2016; Zaldarriaga et al. 2018; Fuller & Ma 2019). Recent work has shown that if post-main-sequence angular momentum transport is not too strong, the first-born BH in BBH progenitors can be highly spinning from either a Case A (main sequence) mass transfer episode or post-main-sequence tidal spin-up (Qin et al. 2019). However, it is unclear if these systems will become BBHs with tight enough orbits to merge within a Hubble time. Alternatively, GW190412 could be of dynamical origin, with the primary BH being the product of one (or more) BBH mergers. The canonical dynamical scenario—formation in a classical globular cluster (Benacquista & Downing 2013)—also struggles to match the parameters of GW190412. To be retained in a globular cluster, the natal spins of first-generation BHs need to be small (e.g., Rodriguez et al. 2019). In this case, the merger product of two BHs will form a second-generation BH with a dimensionless spin of \( \chi \approx 0.67 \): above the measurement of \( \chi_1 \) in GW190412, which is 0.09 \( \leq \chi_1 \leq 0.60 \) with the Phenom approximant and 0.31 \( \leq \chi_1 \leq 0.58 \) with the EOB approximant. The second-generation globular cluster scenario for GW190412’s primary BH is also highly disfavored from phenomenological models of hierarchical mergers, which find an odds ratio of \( \geq 1000:1 \) in favor of a GW190412 being a merger of two first-generation BHs rather than the merger of a first- and second-generation BH in a globular cluster (Kimball et al. 2020). Though globular clusters typically cannot retain higher than second-generation merger products due to the relativistic recoil kicks at merger, nuclear clusters (Gerosa et al. 2020), AGN disks (Tagawa et al. 2020), and high-metallicity super star clusters (Rodriguez et al. 2020) have all been proposed for the formation of GW190412 analogs via hierarchically merging BHs. Other more exotic channels have also been proposed for forming GW190412, such as GW190412 resulting from a 3 + 1 hierarchical quadruple stellar system (Hamers & Safarzadeh 2020), though BBH merger rates from such channels are highly uncertain. Explaining GW observations requires astrophysical models that can produce systems with both parameters and event rates that are consistent with the measured values.

While the formation scenario for GW190412 is to be determined, the correct interpretation of GW190412’s component spins (and those of future GW observations) is paramount for constraining viable formation mechanisms. As the GW detector network continues its observational campaign (Abbott et al. 2018), additional observations of asymmetric and spinning systems (or lack thereof) will further inform the astrophysical channels that lead to the formation of merging BBHs.

Posterior samples for the parameter estimation of GW190412 using our suite of analyses, as well as model evidences, are available on Zenodo (Zevin et al. 2020).

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References
Aasi, J., Abbott, B. P., Abbott, R., et al. 2015, CQGra, 32, 074001
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2018, LRR, 21, 3
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2020, arXiv:2004.08342
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2018, LRR, 21, 3
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2019a, PhRvX, 9, 011019
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2019b, ApJL, 882, L24
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2020, arXiv:2004.08342
Abbott, R., Abbott, T. D., Abraham, S., et al. 2019c, arXiv:1912.11716v1
Acernese, F., Agathos, M., Agatsuma, K., et al. 2015, CQGra, 32, 024001
Ajith, P., Hannam, M., Husa, S., et al. 2011, PhRvL, 106, 241101
Ashton, G., Hübner, M., Lasky, P. D., et al. 2019, ApJS, 241, 27
Ashton, G., & Thrane, E. 2020, arXiv:2006.05039
Babak, S., Taracchini, A., & Buonanno, A. 2017, PhRvD, 95, 024010
Bavera, S. S., Fragos, T., Qin, Y., et al. 2020, A&A, 635, A97
Belczynski, K., Bulik, T., & Fryer, C. L. 2012, arXiv:1208.2422
Belczynski, K., Klencki, J., Fields, C. E., et al. 2020, A&A, 636, A104
Benacquista, M. J., & Downing, J. M. 2013, LRR, 16, 4
Calderón Bustillo, J., Husa, S., Sintes, A. M., & Pürrer, M. 2016, PhRvD, 93, 084019
Chatziioannou, K., Cornish, N., Klein, A., & Yunes, N. 2014, PhRvD, 89, 104023
Chatziioannou, K., Cotesta, R., Ghonge, S., et al. 2019, PhRvD, 100, 104015
Coase, R. H. 1982, How Should Economists Choose? (Chicago, IL: American Enterprise Institute), https://chicagoanbound.uchicago.edu/books/121
Cutler, C., & Flanagan, E. E. 1994, PhRvD, 49, 2658
Damour, T. 2001, PhRvD, 64, 124013
Di Carlo, U. N., Mapelli, M., Giacobbo, N., et al. 2020, arXiv:2004.09525
Eldridge, J. J., Stanway, E. R., Xiao, L., et al. 2017, PASA, 34, e058
Fishbach, M., Farr, W. M., & Holz, D. E. 2020, ApJL, 891, L31
Fishbach, M., Holz, D. E., & Farr, B. 2017, ApJL, 840, L24
Fishbach, M., Holz, D. E., & Farr, W. M. 2018, ApJL, 863, L41
Fragione, G., & Koehs, B. 2018, PhRvL, 121, 161103
Fryer, C. L., Belczynski, K., Wiktorkowicz, G., et al. 2012, ApJ, 749, 91
Fuller, J., & Ma, L. 2019, ApJL, 881, L1
Fuller, J., Piro, A. L., & Jermy, A. S. 2019, MNRAS, 485, 3661
Galaudage, S., Talbot, C., & Thrane, E. 2019, arXiv:1912.09708
Gerosa, D., & Berti, E. 2017, PhRvD, 95, 124046
Gerosa, D., Vitale, S., & Berti, E. 2020, arXiv:2005.04243
Giacobbo, N., & Mapelli, M. 2018, MNRAS, 480, 2011
Giacobbo, N., & Mapelli, M. 2020, ApJ, 891, 141
Graff, P. B., Buonanno, A., & Sathyaprakash, B. S. 2015, PhRvD, 92, 022002
Hamers, A. S., & Safarzadeh, M. 2020, ApJ, 898, 99
Hong, J., Vesperini, E., Askar, A., et al. 2018, MNRAS, 480, 5645
Huang, Y., Haster, C.-J., Vitale, S., et al. 2020, arXiv:2003.04513
Hunter, J. D. 2007, CSE, 9, 99
Kalogera, V. 2000, ApJ, 541, 319
Khan, S., Chatziioannou, K., Hannam, M., & Ohme, F. 2019, PhRvD, 100, 24058
Khan, S., Ohme, F., Chatziioannou, K., & Hannam, M. 2020, PhRvD, 101, 24056
Kimball, C., Talbot, C., Berry, C. P. L., et al. 2020, arXiv:2005.00023
Kimugawa, T., Nakamura, T., & Nakano, H. 2020, arXiv:2005.09795
Klenc, J., Moe, M., Gladysz, W., et al. 2018, A&A, 619, A77
Kushnir, D., Zaldarriaga, M., Kollmeier, J. A., & Waldman, R. 2016, MNRAS, 462, 844
LIGO Scientific Collaboration & Virgo Collaboration 2020, GW190412 data set, doi:10.7935/20yv-ka61
MacKay, D. J. C. 2003, Information Theory, Inference, and Learning Algorithms (Cambridge: Cambridge Univ. Press), https://www.inference.org.uk/itprnn/book.pdf
Mandel, I. 2010, PhRvD, 81, 084029
Mandel, I., Farr, W. M., & Gair, J. R. 2019, MNRAS, 486, 1086
McKinney, W. 2010, in Proc. 9th Python in Science Conf., ed. S. van der Walt & J. Millman, 51, http://pandas.sourceforge.net
Miller, M. C., & Miller, J. M. 2015, PhRv, 548, 1
Miller, S., Callister, T. A., & Farr, W. 2020, ApJ, 895, 128
Nitz, A. H., Dent, T., Davies, G. S., et al. 2020, ApJ, 891, 123
Olejak, A., Belczynski, K., Holz, D. E., et al. 2020, arXiv:2004.11866
Ossokine, S., Buonanno, A., Marsat, S., et al. 2020, arXiv:2004.09421v1
Pan, Y., Buonanno, A., Taracchini, A., et al. 2014, PhRvD, 89, 084006
Pérez, F., & Granger, B. E. 2007, CSE, 9, 21
Posson, E., & Will, C. M. 1995, PhRvD, 52, 848
Puter, M., Hannam, M., & Ohme, F. 2016, PhRvD, 93, 084042
Qin, Y., Fragos, T., Meynet, G., et al. 2018, A&A, 616, A28
Qin, Y., Marchant, P., Fragos, T., Meynet, G., & Kalogera, V. 2019, ApJL, 870, L18
Racine, E. 2008, PhRvD, 78, 044021
Rodriguez, C. L., Kremer, K., Gradic, M. Y., et al. 2020, ApJL, 896, L10
Rodriguez, C. L., & Loeb, A. 2018, ApJL, 866, L5
Rodriguez, C. L., Zevin, M., Amaro-Soane, P., et al. 2019, PhRvD, 100, 43027
Rodriguez, C. L., Zevin, M., Pankow, C., Kalogera, V., & Rasio, F. A. 2016, ApJL, 832, L3
Roulet, J., & Zaldarriaga, M. 2018, MNRAS, 478, 4216
Safarzadeh, M., & Hotokezaka, K. 2020, ApJL, 897, L7
Santamaria, L., Ohme, F., Ajith, P., et al. 2010, PhRvD, 82, 064016
Schmidt, P., Ohme, F., & Hannam, M. 2015, PhRvD, 91, 024043
Smith, R., Ashton, G., Vajpeyi, A., & Talbot, C. 2019, arXiv:1909.11873
Tagawa, H., Haiman, Z., Bartos, I., & Kocsis, B. 2020, arXiv:2004.1914
Talbot, C., & Thrane, E. 2017, PhRvD, 96, 023012
Thrane, E., & Talbot, C. 2019, PASA, 36, e010
Van Den Broeck, C., & Sengupta, A. S. 2007, CQGra, 24, 1089
Van Der Walt, S., Colbert, S. C., & Varoquaux, G. 2011, CSE, 13, 22
Wright, T., & Ajith, P. 2017, PhRvD, 96, 124024
Zackay, B., Venumadhav, T., Zackay, B., Roule, J., Dai, L., & Zaldarriaga, M. 2020, PhRvD, 101, 83030
Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, NaMe, 17, 261
Wysocki, D., Lange, J., & O'Shaughnessy, R. 2019, PhRvD, 100, 043012
Zackay, B., Venumadhav, T., Dai, L., Roulet, J., & Zaldarriaga, M. 2019, PhRvD, 100, 23007
