Mechanics of metal-catecholate complexes: The roles of coordination state and metal types

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There have been growing evidences for the critical roles of metal-coordination complexes in defining structural and mechanical properties of unmineralized biological materials, including hardness, toughness, and abrasion resistance. Their dynamic (e.g. pH-responsive, self-healable, reversible) properties inspire promising applications of synthetic materials following this concept. However, mechanics of these coordination crosslinks, which lays the ground for predictive and rational material design, has not yet been well addressed. Here we present a first-principles study of representative coordination complexes between metals and catechols. The results show that these crosslinks offer stiffness and strength near a covalent bond, which strongly depend on the coordination state and type of metals. This dependence is discussed by analyzing the nature of bonding between metals and catechols. The responsive mechanics of metal-coordination is further mapped from the single-molecule level to a networked material. The results presented here provide fundamental understanding and principles for material selection in metal-coordination-based applications.

Many biological materials with outstanding mechanical, adhesive and frictional properties benefit from their hierarchical composite structures, within which the reinforcing role of metal complexation in the polymer network phase was recently highlighted1-2. For example, it was reported that the Fe-catecholate (1,2-dihydroxybezene) complexes in mussel byssus cuticles provide effective crosslinking strategies between metal ions and catecholic amino acid dopa (3,4-dihydroxypheylalanine) in mussel foot protein 1 (mfp-1) chains3-5. Granules with high density of these crosslinks resemble the particle-reinforced phase in the composite, and has correlated to the hardness and high extensibility of the mussel byssus threads. The metal-dopa crosslinks in the matrix were also proposed to endow self-healing properties. Measurements of the mechanical properties after removing metals from the cuticle showed 2-fold decrease in the hardness, suggesting a significant impact of the coordination complex on the material properties3. In the fangs of the wandering spider Cupiennius salei, Zn-His coordination complexes were also found in the matrix of chitin and Zn-containing proteins6. Connection was then identified between the mechanical properties and Zn/Cl enrichment in this fiber-reinforced biocomposite. Moreover, the inhomogeneous metal distribution in the fang might have additional implications in its graded mechanical properties, with an over two-fold increase in hardness and 1.5-fold increase in stiffness measured from the base to the tip6.

The aforementioned examples give prominence to the key roles of metal coordination in defining mechanical properties of crosslinked biopolymers. Coordination complexes form crosslinks when two or more ligands each donate a nonbonding electron pair to empty orbitals in a transition metal ion. A hexadentate mononuclear triple (tris-)Fe-dopa coordination complex was proposed for the mfp-1 crosslinks in the byssus coating3, while bi-, quadra- and hexa-dentate complex can be induced in synthetic materials by tuning pH4. In comparison with other crosslinking strategies in polymeric networks, such as non-bonding (e.g. van der Waals, electrostatic, hydrogen bonds)5 and covalent bonds6, metal-coordination-based crosslinks play a unique role as they are strong (comparable to covalent bonds), responsive to the environment (e.g. pH), and yet reversible (to prevent catastrophic material failure)7-9. These features have been proposed to offer biological materials a number of desirable mechanical properties in the absence of mineralization, including enhanced hardness, adhesion, toughness, tunability, self-healing capability, efficient energy dissipation upon impact loading, and functions as sacrificial bonds10-12.
Recently efforts have been not only made to understand biological materials, but also to develop synthetic materials following these design strategies (Fig. 1). Possible modifications of the proteins and their interfaces were pursued to establish demanding mechanical properties. Metal-coordination hydrogels were successfully synthesized by several groups, which display distinctly dynamic (e.g. responsive, self-healable, tunable, etc.) material properties in contrast to conventional polymers. One excellent example is the pH-responsive gel inspired by the pH jump that proteins experience during maturation of a mussel byssus secretion, which features high stiffness approaching covalently crosslinked gels with additional self-healing capability. In order to develop and optimize bio-inspired materials following the recipes found in nature and explore the limit of maximizing material performance, the understanding of metal-coordination-based crosslinks, especially their mechanical significance, is urged.

To address these issues, we present here first a molecule-level study of the coordination complexes. First-principles calculations were performed to investigate representative coordinate and covalent crosslinks, with focus on their structural and mechanical properties. The impact of these crosslinks to mechanics of materials at the network level is then explored, which reveals the responsive hallmark of metal-coordination based biological and bio-inspired materials.

Results
We considered both metal-catecholate coordination complexes and covalent bonds as possible crosslinking mechanisms in biopolymer networks. Ultraviolet-visible (UV) absorption spectra studies showed that at low pH (<7), catechol is reduced and coordination complexes are dissolved instead of crosslinked. At higher pH, double (bis-) and triple (tris-)complexes form. On the other hand, covalent bonds could form by an oxidative process under appropriate (pH 8.5) conditions. Our first-principles calculations start from atomic structures illustrated in Fig. 1c. In the bis-complexes, a planar square coordination complex was assumed, while the metal ion locates in the center of an octahedron for tris-complexes.

Five metals, including Fe, Mn, Ti, Zn, and Ca, were included in our calculations according to their wide appearance in biological (Fe, Mn, Zn, Ca) and synthetic (Fe, Ti) polymer materials. Our geometry optimization results show that all metals form distorted planar square and octahedron configurations with their oxygen donors provided by surrounding catechols. The lengths of metal-oxygen (MO) bonds in the bis-complexes were found to be 1.85 (Fe-O), 1.88 (Mn-O), 1.92 (Ti-O), 2.03 (Zn-O), 2.01 (Ca-O) Å in the bis-complexes, and 1.88 (Fe-O), 1.90 (Mn-O), 2.00 (Ti-O), 2.13 (Zn-O), 2.38 (Ca-O) Å in the tris-complexes, respectively. The O-M-O bond angles are distorted from 90° (the angle in undistorted planar squares and octahedrons) for both bis- and tris-complexes. The lengths of metal-oxygen (MO) bonds in the tris-complexes were found to be 1.85 (Fe-O), 1.88 (Mn-O), 1.92 (Ti-O), 2.03 (Zn-O), 2.01 (Ca-O) Å in the bis-complexes, and 1.88 (Fe-O), 1.90 (Mn-O), 2.00 (Ti-O), 2.13 (Zn-O), 2.38 (Ca-O) Å in the tris-complexes, respectively. The O-M-O bond angles are distorted from 90° (the angle in undistorted planar squares and octahedrons) for both bis- and tris-complexes.
catechol groups was also explored, suggesting a length of 1.49 Å for the inter-catechol C-C bond.

The coordinate bonds are commonly considered to be comparable to, but weaker than covalent bonds. When the polymer network are loaded mechanically as mimicked in this work by pulling apart catechol groups, the metal ion-catechol crosslinks bear tensile loads. The coordination complex thus not only critically defines the stiffness, strength, viscoelasticity, and other overall mechanical behaviors, but also being responsible for the reversibility and tunability of the networked material. Here we quantified the mechanics of crosslinks by displacing and constraining selected carbon atoms in catechols (see Methods section for more details). The force carried by the complex was calculated from the reaction force on the loaded or fixed atoms, and the results are plotted as a function of relative displacement (defined as change in the distance between pulled and fixed carbon atoms) in Fig. 2. It was found that in comparison to the covalent crosslink, the mechanical response of coordination complexes feature a rather complex profile with multiple peaks, arising form a larger number of degree of freedom involved in the deformation (bond stretching, bond angle bending, etc.). For bis-complex, the ridge was first flattened with negligible responses in the force amplitude. Further stretching led to subsequent breaking of MO bonds (see Fig. 3 and Movie M1 and M2 in Supplementary Information). The irregular profiles of force-displacement curves in Fig. 2, especially

![Graphs showing force-displacement relationships for metal (Fe, Mn, Ti, Zn, Ca)-catechol complexes, and (f) for the covalent crosslink. Metal-coordination shows tensile strength of the complex on the order of a few nN, varying by the type of metals.](image-url)
those for tris-complexes, can be explained from the dynamics of MO bond distortion and detailed breaking sequence. The Fe-catechol bis-complex first experienced a structural distortion, corresponding to the first peak, and then failed by breaking two Fe-O bonds simultaneously, while for Ti-catechol bis-complex, Ti-O bonds on the pulled catechols broke first. The tris-complexes also showed different bond breaking patterns under tensile loading. For Fe-catechol tris-complex, the two pairs of metal-catechol crosslinks at the left side of the metal ion (Fig. 1c) broke in two steps, resulting in failure at a much larger distance than the Ti-catechol tris-complex, where the breakage of crosslinks right to the metal ion led directly to structural failure (Fig. 3). These distinct behaviors suggest different stabilities of metal-coordination complexes at nonequilibrium conditions, although their structures at equilibrium are quite similar as described above.

To further assess the roles of metal ions and their coordination states, the whole spectrum of tensile strength for metal-coordination complexes was calculated and summarized in Fig. 4. For bis-complexes, the strength defined as the peak force carried by the crosslink consisting of two MO bonds ranges widely from 3.30 (Ca-O), to 8.18 (Ti-O) nN. Notably, the Fe-catecholate bis-complex features a tensile strength of 6.22 nN, while Ti-coordinate crosslink has a tensile strength of 8.18 nN. It should be noted here that this strength defined for the coordination complex is actually contributed by a pair of M-O bonds and thus the metal coordination could reach up to half of the strength in a covalent bond.

Figure 3 | Structural distortion and failure of coordination complexes. We plot here changes in the metal-oxygen bond lengths (labeled in the insets) in the metal-coordination complexes with respect to the loading distance (a, c: Ti, b, d: Fe).

Figure 4 | Dependence of crosslinking strength on coordination state and metal types. Tensile strength of coordination complexes based on different metals (Fe, Mn, Ti, Zn, Ca) and catechols are plotted and compared with a covalent crosslink between two catechol groups. It should noted that the strength of complexes are contributed by two metal-oxygen bonds and thus the metal coordination could reach up to half of the strength in a covalent bond.
The dependence of tensile strength on both the coordination state and metal ion types shares the same trend as bis-complexes, i.e. Ti-based coordination complex are the strongest, followed by Fe and Mn featuring similar performance, while Zn and Ca show lower values of the strength. Moreover, we could extract the tensile stiffness from the linearly elastic part of the curve (Fig. S1). The results show that for bis-complexes, Fe-based and Ti-based complex could offer stiffness \( k \) of 75.61 and 78.28 N/m respectively, also close to the covalent crosslink with \( k = 75.77 \) N/m, which is one order higher than the hydrogen bonds. Similarly, the stiffness of single M-O bond is half of the stiffness measured for the coordination complex.

The structural and mechanical properties of the coordination complexes can be clarified further from their electronic structures. Bader atomic charge analysis was performed to inspect charge distribution in the complexes. This approach divides the grid-based charge density from quantum mechanical calculations of electrons into atomic volumes. The analysis shows that for bis-complexes, 1.3, 1.7, 1.9, 1.3, and 1.4 electrons are transferred from the metals (Fe, Mn, Ti, Zn, Ca) to the catechols. For tris-complexes, the amount of electrons transferred is 1.3, 1.4, 1.8, 1.3, and 1.5 electrons respectively. These results suggest more ionic nature for divalent metal such as Ca, which is correlated to the mechanical strength as identified above. These results elucidate the combined electrostatic and covalent nature of the metal-coordination. To quantify the difference in the metal-catechol bonding characteristics, topological analysis based on electron localization functions (ELFs) was performed for all complexes. ELF is a localized function of the ground-state electron density and wavefunction obtained from the first-principles calculations. The value of ELF ranges from 0 to 1, where 1 corresponds to the perfect localization as in covalent bonds, and 0.5 corresponds to the electron-gas-like pair probability as in metallic bonds. The bonding nature between metal ions and oxygen atoms in catechols depends on the type of metal. As shown in Fig. 5a and 5b, the results illustrate the nature of coordination bonding compared to the covalent and ionic bonds. Moreover, there is a distinct contrast in the bonding nature between Ti-O and Ca-O bonds. Further comparative analysis shows that only Ca in the five metals explored here displays distinctly the ionic nature in the coordination complex.

The superior mechanical performance revealed above for single metal-catecholate complex could inspire synthetic materials by incorporating metal coordination. However, how could these merits of individual crosslinks be mapped to the material level remains as a critical question to answer. Especially, the contrastive stiffness of the complex to non-bonding interactions opens promising opportunities in designing highly responsive materials. We now discuss this point using a crosslinked polymer network mode with the assumption of affine deformation. The mechanical response of a typical crosslinked polymeric network can be described using a few representative units (Fig. 6a). When the representative units are loaded, both the polymer chain and crosslink bears an internal force \( f \). We define the effective tensile stiffness of polymers (considering their intrinsic elasticity, thermal fluctuation, and network entanglement) to be \( K \), and the stiffness of crosslinks to be \( k \) respectively. At its mechanical equilibration, the nature of coordinate bonds could be modified into a new state with stiffness \( k' \), e.g by tuning the pH, type or concentration of metal ions. Then a response in the force amplitude is generated inside the material. Straightforward analysis shows that response is \( f' = f(a + b)/b(a + b) \), where \( a = k'/k \) is the stiffening factor of crosslinks, and \( b = K/k \) is the relative stiffness between the polymer and crosslink. Typical scaling behavior of \( f'/f \) with dependence on the stiffness enhancement \( b \) is plotted in Fig. 6b. The results show that significant response can be established in a stressed network, which increases with both \( a \) and \( b \). For \( b = 10 \), a ten-fold stiffening of the crosslink \((a = 10, e.g. from hydrogen bonding to metal-coordination) could yield nearly ten-fold response in the internal force carried by the material. We also find that the response is more amplified in the softening situation, i.e. with \( a < 1 \).

**Discussion**

It was shown in the results that the highest strength achieved in Ti-catecholate complex could reach half of the covalent strength. Recently, atomic force microscopy (AFM) measurements suggest that the strength of a covalent bond is ~2.0 nN (for silicon-carbon) and ~1.4 nN (gold-sulfur). The mechanics of metal coordination in biological materials was characterized as well in AFM and surface force apparatus (SFA) studies. The results show that a single dopa

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**Figure 5** | Constructive nature of metal (Ti, Ca)-catecholate bonding in the coordination complex. Isosurfaces at ELF = 0.5 are plotted here with the atomic structures of tris-complexes. The distribution of ELF shows distinctly covalent and ionic nature for (a) titanium and (b) calcium respectively, which explains their difference in the mechanical properties (Fig. 2).

**Figure 6** | Responsive mechanics of crosslinked networks. (a) Typical crosslinking structures in a stressed polymer network can be captured in a simple spring model, loaded by force \( f \). (b) Response in the internal forces \((f'/f)\) carried by a pre-stressed networks (the inset), with dependence on the change in tensile stiffness \((k'/k)\) of coordinate crosslinks. \( k \) is the effective stiffness of the polymer, and \( k \) is the stiffness of the crosslink.
residue contacting a wet TiO$_2$ surface reveals high strength (~0.8 nN) yet fully reversible, noncovalent interaction. Moreover, Fe ions can mediate the unusually strong interaction between positively charged proteins through metal coordination with dopa. Although there are some theoretical and experimental evidences showing that mono- and bis-dentate coordination could form between catechol and a TiO$_2$ surface, as a result of dissociative adsorption, the detailed single-molecule-level chemistry at interfaces in the AFM and SFA experiments is not known. Moreover, the interfacial interaction and stability of metal-coordination complex could be perturbed by its chemistry and environment such as the presence of water. In spite of these uncertainties that prevent rigorously quantitative comparison, our results that are on the same order as the reported data remain to be conclusive.

To further explain the quantitative discrepancy between our calculations and experimental measurements, it should be noted that the failure of metal coordination complexes has dissociation rates over a range of time scales. Both thermal fluctuation and loading rate will strongly affect the bond dissociation. These effects can be essentially captured by the Bell-Evans model $f_d = k_B T n \left( \tau_0 \right) / k_0$, where $k_B$ is the Boltzmann factor, $T$ is the temperature, $\tau_0$ is the loading rate, $\tau_0$ is the equilibrium bond length, and $k_0$ is the equilibrium dissociation rate. Recent experiments have shown the evidence that the stiffness and strength of synthetic gels containing metal-coordination crosslinks are rather a function of the dissociation rate than the equilibrium constant. Our quasi-static loading process thus approaches the limit of $r = 0$ and $T = 0$. In the force microscopy tests, although finite loading rate would yield a higher strength than that obtained in our calculations, the thermal fluctuation is expected to play a more critical role, which significantly reduces the lifetime of metal-ligand bonds and thus the measured strength of coordination complexes. This may explain the relatively low values reported in experiments. Additionally, although specific metal-coordination complex was modeled here according to experimental evidence, explicit effects from the chemical environment (pH, ionic strength, salt, etc.) could be critical. For example, water or hydroxyl could fill in the empty octahedral sites in the bis-complex, ions and molecules around could induce charge transfer and polarization as well. The solvent effect for metal ions in aqueous solution could vary for different metals, and thus lead to further contrastive mechanical properties of the complexes. However, due to the high computational demands to include these effects in first-principles calculations, we leave them to be exploited in our future work.

The nature of coordinate bonds is predominantly electrostatic and partially covalent, which differs for different combination of metal and donor atoms. There are some understandings from a survey of metals in protein structures, that the covalent contributions follow a certain pattern: $K^+ < Na^+ < Ca^{2+} < Mn^{2+} < Fe^{2+} < Co^{2+} < Ni^{2+} < Cu^{2+} < Zn^{2+}$. In comparison to the metal-donor bond lengths between metal and oxygen in main-chain carbonyl reported in Ref. 31, the atomic structures of catechols and typical coordination complexes in metal oxides were chosen as our model systems here, because catechol and its derivatives are widely present in biological materials, and contained in the side chains of dopa. Also catechol serves as a minimal model to reduce the computation. Supercells with dimensions of 20, 16, and 16 Å were built. The thickness of vacuum layers was verified to be large enough to model the complexes as isolated molecular systems, by comparing with results using larger supercells.

First-principles calculations. The structure and properties of the coordination complexes were explored using plane-wave-basis-set-based density functional theory (DFT) methods. Spin-polarization was included in the calculations. Generalized gradient approximation (GGA) in Perdew–Burke–Ernzerhof (PBE) parameter settings was used for the exchange–correlation functional and Projector Augmented Wave (PAW) potentials were used for ion–electron interaction. We used the Vienna Ab-initio Simulation Package (VASP) for the calculations. For all results presented, an energy cut-off of 500 eV was used for the plane-wave basis sets. Gamma point was used for the Brillouin zone integration. These settings were verified to achieve a total energy convergence less than 1 meV per atom. For geometry relaxation, the force on atoms converged below 0.01 eV/Å. All structures were optimized using the conjugated gradient method. This first-principles approach was successfully applied to similar systems, e.g. adhesion of catechol on TiO$_2$ and silica. In previous studies, adhesion energy close to experimental measurement was predicted. However, it should also be noted that the high computational demand prevents DFT calculations to include explicit solvent, and thus the effects of chemical environment (e.g. pH, ionic strength, salt, etc.) could not be directly addressed. In this study, this information was fed into the molecular models instead, as suggested by absorption spectra studies of metal-catecholate complexes. Optimizing the models in first-principles calculations could eventually solve this issue by minimizing computational
demands while keeping the essential chemistry. This is the subject of our ongoing investigation.

Mechanical tensile tests. Tensile loads were applied on selected carbon atoms in the catechols. For the bis–catecholate- and tris–catecholate-complexes, the rightmost two carbon atoms (referring to the metal ion, see Fig. 1c) were displaced further from the other catechols step by step (0.1 Å for each step), while the leftmost two carbon atoms (one in each catechol group) were fixed in the loading direction while not constrained in other directions. Coords of all atoms in the system were scaled accordingly after each loading step, to minimize the dependence on the step size. Geometrical optimization calculations were performed after each movement, and the force was evaluated from the reaction force on displaced or constrained atoms. For the covalent bonds between two catechols, the same loading procedure was followed except that displacement and constraints were only applied to one atom in each catechol group.

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