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Consider the simplest cells, such as bacteria or yeast. Cells grow at different rates, depending on their environment. A cell's growth rate depends on how much food is present, on the temperature and salt concentration of the external medium, and on its internal biochemical health. Because a cell's duplication speed is often the single most important determinant of its ability to propagate its progeny, growth rate could have evolved to be a complicated function of many biochemical details of a cell. However, we review here recent efforts toward a different view. Modeling shows how the growth laws of simple cells are encoded within the physical properties of a cell's proteome (i.e., its full complement of proteins). That is, some cell behaviors are attributable to large fractions of the proteome, not just a single protein or gene or pathway. And, some behaviors are physical (due to protein folding, aggregation, or diffusion, applicable in some universal or general way across different proteins), rather than biological (due to the protein's particular biological action). Of course, at best, simple models of the physical proteome are only a first approximation. But, in the spirit of other physical chemistry, they may provide useful conceptual insights and can make testable predictions.

First, we make a general point: growth laws are related to, and manifestations of, evolutionary fitness landscapes. Define a cellular growth rate, $\lambda$, as the number of new cells produced per unit time from each existing parent. If $c(t)$ is the cell population at time $t$, then under appropriate conditions, populations grow as

$$\frac{dc}{dt} = \lambda c$$

The growth rate $\lambda$ can depend, often strongly, on various quantities; these are called growth laws. Perhaps the best known growth law, $\lambda = \lambda(sugar)$, indicates that cells grow faster with increasing concentrations of food, such as sugar, up to a point at which the growth rate saturates. Bacterial growth rates also depend strongly on temperature and external salt concentrations. For practical bacteriology, these are important. To kill bacteria, you remove a food source, or you heat the cells to high temperatures (as when you cook food), or you introduce high external salt concentrations (in pickling fish or in making jerky or salting meats, for example). In general, such growth laws can be expressed as $\lambda = \lambda(e)$, where $e$ indicates a vector of environmental variables, such as sugar, temperature, or salt. These functions can express cellular growth laws.

A growth law is a function that describes how "today's cell" can respond to variations in today's conditions. But, cells can change those functions, through evolutionary modifications over longer time scales. This can be expressed in terms of their genotype, a vector of genes, $g$. We use the term genotype here in a very general way: It can describe either a set of discrete options, such as the presence or absence of genes or amino acids in proteins, or a continuum of options. It can express some property of a gene directly or it can be a surrogate for that, representing some rate coefficients or equilibrium constants in the biochemical workings of the cell. In general, we can express the growth rate of a cell as

$$\lambda = \lambda(e, g)$$

Role of Proteome Physical Chemistry in Cell Behavior

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ABSTRACT: We review how major cell behaviors, such as bacterial growth laws, are derived from the physical chemistry of the cell's proteins. On one hand, cell actions depend on the individual biological functionalities of their many genes and proteins. On the other hand, the common physics among proteins can be as important as the unique biology that distinguishes them. For example, bacterial growth rates depend strongly on temperature. This dependence can be explained by the folding stabilities across a cell's proteome. Such modeling explains how thermophilic and mesophilic organisms differ, and how oxidative damage of highly charged proteins can lead to unfolding and aggregation in aging cells. Cells have characteristic time scales. For example, $E. coli$ can duplicate as fast as 2–3 times per hour. These time scales can be explained by protein dynamics (the rates of synthesis and degradation, folding, and diffusional transport). It rationalizes how bacterial growth is slowed down by added salt. In the same way that the behaviors of inanimate materials can be expressed in terms of the statistical distributions of atoms and molecules, some cell behaviors can be expressed in terms of distributions of protein properties, giving insights into the microscopic basis of growth laws in simple cells.
Equation 2 captures both today’s growth law $\lambda = \lambda(\mathbf{g})$, for fixed evolutionary properties $\mathbf{g}$, while it also captures that growth rates can be modulated by evolution $\dot{\lambda} = \dot{\lambda}(\mathbf{g})$ for fixed conditions $\mathbf{e}$. The latter property, $\dot{\lambda} = \dot{\lambda}(\mathbf{g})$, is the fitness landscape for cells for which duplication speed is their primary measure of fitness. Hence, Eq 2 relates, albeit in only a general abstract way, the evolutionary fitness landscape to the growth laws of cells. For cells that have been under a fixed selection pressure for a long time, and have evolved to maximize their fitness, we can study their peak-fitness points by finding

$$
\frac{d\lambda}{dg} \bigg|_{g^*} = 0
$$

Note that, in general, cellular fitness $f$ is not always equal to just $\lambda$, the growth rate. Many types of cells live in multicellular organisms. They contribute to the fitness of the whole organism. Their own particular fitness objectives are rarely known. Here, we describe some models of fitness $f(\mathbf{e}, \mathbf{g})$ in simple cells as a function of properties of the cell’s proteins.

We focus on proteins because more than half of a cell’s biomass is its proteins. Hence, where physical behaviors matter, proteins are likely to be predominant players. We distinguish between a protein’s generic physicochemical properties and its specialized sequence-structure actions. By “general physical” properties, we mean the following. First, we are referring to a protein’s health (also called proteostasis): the balance between folded and unfolded states, the balance between folding and degradation, and the states of protein oxidation. Second, we are also referring to biophysical properties that can matter to the cell, such as protein movement, transport, crowding, sticking, and localization. Thanks to enzymatic assays, genome sequencing, and tens of thousands of atomically detailed protein structures in the Protein DataBank, the special functions of many proteins are now known. Less is known about the generic, physical, and health behaviors of proteomes. While the biological actions are often distinct from one protein to the next, the physical behaviors can involve commonalities among proteins, often arising more from statistical properties than from the singular native states. These properties include a proteome’s distribution of stabilities, folding rates, and sensitivities to perturbations (such as side-chain charge modification), as shown in Figure 1. The physical properties of proteins are important because the cell commits major resources in energy and biomass toward managing them, in its struggle against stresses, disease, and death. Just like the specialized jobs of proteins, the generic actions can be changed through evolutionary processes such as natural selection.

Here, we describe how simple physicochemical models, combined with data from in vitro experiments, can predict some cell behaviors, rationalize observed growth laws, and generate hypotheses about diseases, aging, and evolutionary tendencies. The concepts being sought here, and the models being developed, are coarse-grained, not atomically detailed. Yet, despite their simplicity, they are often sufficient to generate testable hypotheses. The first example below shows how a coarse-grained model of protein folding stability can explain the high sensitivities of cells to temperature, rationalize thermal growth laws, predict proteome stability distribution functions, and give insight into how thermophilic organisms may have evolved to deal with higher environmental temperatures.

### THERMAL PROPERTIES OF CELLS ARISE FROM THE FOLDING STABILITIES OF THEIR PROTEOMES

Cells are highly sensitive to temperature. It is not uncommon that the temperatures at which cells die are only a few degrees higher than the temperatures at which their growth is optimal. Small shifts of environmental temperature can drive biological migrations, extinctions, genetic divergence, and speciation. By what mechanism are cells so sensitive to temperature? Here, we review a polymer folding model (polymer-collapse theory) that indicates that the thermal sensitivities of cells arise because proteomes have evolved to have denaturation temperatures that are only marginally higher than the cell’s growth temperature. Despite its simplicity, this mechanism gives an approximate quantitative description of bacterial growth rates versus temperature.

**Cells Are Sensitive to Temperature Because Their Proteomes Are Poised Near Their Denaturation Temperatures.** This protein—denaturation—catastrophe mechanism has been made quantitative by a combination of thermodynamic measurements of 59 mesophilic proteins in vitro with polymer-collapse theory. Such theory reckons that reversible protein folding is driven by the small average tendency of amino acids to prefer sticking to other amino acids inside a compact native structure, rather than to be exposed and solvated in an expanded unfolded state in water. This mechanism reckons that the principal force opposing folding is the chain entropy, which favors the unfolded state. A version of that simple idea also accounts for electrostatic interactions.
and the effects of temperature, salts, and denaturants, giving the folding free-energy $\Delta G_{\text{unfold}} = G_{\text{unfolded}} - G_{\text{folded}}$ as:

$$
\frac{\Delta G_{\text{unfold}}(T, \text{pH}, c_{e}, c)}{kT_0} = -N \left[ \frac{g_0 + m_{i}c}{kT_0} + \frac{\Delta c}{kT_0} (T - T_0) + \frac{T}{T_0} \frac{\Delta c}{T} \ln \frac{\Delta c}{T} + \frac{k \Delta c}{2N} \ln 2(1 + xR_0) - \frac{Q_2}{R_0(1 + kR_0)} \right],
$$

(4)

where $g_0$ represents the free-energy when amino acids desolvate and come into contact, $z$ is the average conformational freedom loss per backbone bond, and $\Delta c$ is the change in heat capacity per amino acid upon folding. $Q_4$ and $Q_6$ are the total net charge on the denatured and native structures, respectively, and $R_4$ and $R_6$ are the radii of denatured and native protein. $N$ denotes the number of amino acids (or chain length) in the protein, $c$ is the denaturant concentration, $k$ is the inverse Debye length, $\kappa$ is Boltzmann’s constant, $T$ is the temperature, $T_0 = 373.5$, and $T_s = 385$ K. For details, see refs 10 and 14.

Equation 4 gives the stability for a single average protein of length $N$. Thus, the probability distribution $p(\Delta G)$ of stabilities of all the proteins in a proteome (Figure 1a) can be computed from $P(N)$, the distribution of chain lengths of proteins in a cell. $P(N)$ is available for different cell types from proteomic or genomic data.

We conclude that proteomes tend to be marginally stable at their physiological temperatures; see Figure 2. This marginal stability is not because the average stability is low, but because of the distribution of stabilities. The average protein in E. coli is estimated to be reasonably stable, $\Delta G_{\text{unfold}} = 6.8$ kcal/mol at 37 °C. However, there are many proteins that populate the “unstable” side of the distribution: approximately 550 out of 4300 (size of the E. coli proteome) proteins are less stable than 3 kcal/mol. In the absence of much data, we can estimate how stability is affected by protein domain structure, and it indicates that proteins may be even less stable than the estimates above. Furthermore, while these estimates are based on stabilities measured in vitro, experiments and simulations show that protein stabilities in vivo or in the reconstituted cytosol are comparable to, or even slightly less stable than, those in vitro. The polymer folding model predicts that this marginally stable subset of the proteome is responsible for the high thermal sensitivity of the cell, as seen in Figure 2 by a small shift in temperature from 37 to 41 °C.

A similar stability distribution is predicted by an evolutionary kinetics model. In that treatment, random mutations occur through evolution that can alter the folding stabilities of proteins. Evolutionary changes occur by a random walk with a drift on the folding free-energy landscape. That work envisions two limiting states. Proteins have a maximum stability, $\Delta G_{\text{max}}$, because it becomes increasingly harder for evolution to find sequences having arbitrarily high stabilities. Proteins also have a minimum stability, $\Delta G_{\text{min}}$ because otherwise they will aggregate or not fold. Within these two limits, it is assumed that the fitness landscape is flat. The protein stability distribution that evolves through this evolutionary model gives the same stability distribution as the polymer folding model.

Both the polymer folding model and the evolutionary kinetics model give a basis for rationalizing the functional form of cellular thermal growth laws. We suppose that the cell’s growth rate, $r(T)$, is a product of two terms: (i) a factor that describes Arrhenius-activation of one or more activated metabolic process(es) that govern how the cell’s growth rate increases with temperature at low temperatures, and (ii) a factor that accounts for the fraction of the proteome that is folded at any temperature (capturing the denaturation catastrophe of the proteome at high temperatures):

$$
r(T) = r_{0} \exp \left( \frac{-\Delta H^\ddagger}{kT} \right) \prod_{i=1}^{\Gamma} \frac{1}{1 + \exp(-\Delta G_{\text{unfold}}(N_i, T)/RT)}
$$

(5)

Here, $r_0$ is some reference growth rate, $\Delta H^\ddagger$ is the activation barrier of some critical growth-limited metabolic rate, and $\Gamma$ is the number of essential proteins that are needed for growth. The product denotes multiplication over the probability that the $i$th essential protein (with $N_i$ amino acids) is in the folded state which is written in terms of $\Delta G_{\text{unfold}}$ (eq 4; typical temperature dependence shown in Figure 3a). The expression above is simplified by assuming lethal proteins are drawn from the same distribution as the proteome, thus enabling the calculation over all the proteins in the proteome, with $\Gamma$ being a fit parameter. The details of the calculation can be found in previous work. Similar arguments have been made but using only a single effective value for $\Delta G_{\text{unfold}}$. The model described here, based on the whole proteome stability distribution, fits well the experimentally measured growth rates for mesophilic organisms (Figure 3b). The corresponding best-fit value of the cell’s activation barrier for growth, $\Delta H^\ddagger$, for E. coli is found to be 16.3 kcal/mol. This happens to be approximately equal to the barrier for peptide bond formation by the ribosome, and is consistent with estimates from other studies. Moreover, this activation energy is in the same range as typical values for various enzymatic reactions, including the barrier (13 kcal/mol) that is associated with the elongation of RNA by transcription. This model also fits the growth rates of thermophilic organisms (Figure 3c) well when

![Figure 2. Distribution of unfolding free-energy ($\Delta G_{\text{unfold}} = G_{\text{unfolded}} - G_{\text{folded}}$) of all the proteins present in the E. coli proteome at 37 °C (blue) and at 41 °C (red). The bin width for the free-energy is 1 K. The total area under the curve equals the number (4300) of proteins present in the E. coli proteome. Adapted with permission from ref 8. Copyright 2010 Elsevier.](image-url)
using thermodynamic parameters for thermophilic proteins obtained from analyzing in vitro data sets.10 A detailed systems level model has been applied to understand how mutations in metabolic networks change thermal growth rates.27,28 They also indicate that the thermostabilities of metabolic enzymes are rate-limiting at superoptimal temperatures.28 These models and arguments suggest that fundamental physicochemical properties of proteomes help to define a cell’s evolutionary fitness landscape (Figure 3d).

Proteomes of Thermophilic Organisms Are More Stable Than Those of Mesophilic Organisms. The polymer-collapse model also gives insight into how mesophilic cells differ from thermophiles. Mesophilic organisms mostly live at moderate temperatures (25−40 °C) while thermophilic organisms grow at higher temperatures. How do their proteomes differ? A global analysis of 57 thermophilic proteins and 59 mesophilic proteins shows an average systematic difference:10 thermophilic proteins denature at higher temperatures than mesophilic proteins, as they are more stable, on average, at all temperatures10 (see Figure 3a). It also indicates that denatured states of thermophilic proteins may have less chain entropy than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This suggests that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10 This implies that the denatured states are, on average, more compact than mesophilic proteins.10

However, it seems likely that electrostatics may be a key contributor to these differences.39−51 Electrostatic stability of folded proteins can depend both on a protein’s net charge and on its charge patterning. For example, Sawle and Ghosh have shown that a good predictor of the relative compactness of the denatured structures between thermophic and mesophilic sequences is the sequence−charge−decoration (SCD) metric57

\[
SCD = \frac{1}{N} \sum_{m=2}^{N} \sum_{n=1}^{m-1} q_m q_n (m - n)^{1/2}
\]

Here, \(q_m\) and \(q_n\) are the charges (1 for basic, −1 for acidic, and 0 otherwise) on two amino acids \(m\) and \(n\) with \(lm - nl\) being their sequence separation. SCD expresses the degree of charge mixing; a similar metric has been given by Das and Pappu.55 Figure 5 gives the SCD values for two sequences of charge. A more compact denatured state is predicted by a more negative value of SCD. In this case, a “blockier” sequence of charges gives the more compact denatured state. Sawle and Ghosh have applied this metric to a set of 540 orthologous pairs of thermophilic and mesophilic proteins, and found that thermophiles, in general, have a more compact denatured state than mesophiles.57 While this comparison was made without corresponding 3D protein structures, a comparison has also been made of a smaller set of 55 well-aligned mesophile−thermophile pairs, for which structures are known.56 This too
some proteins be sufficient to contribute to the aging phenotype?

Here, we review the following mechanism:72 (i) oxidation damages amino acid sites on random proteins across the proteome; (ii) some damage events will alter the charges on some side-chains;77 (iii) for a small subset of the proteome, a small change in net charge (as small as +1 or −1 charges) can denature or destabilize its folded state. How can changing a protein’s charge by only +1 or −1 units unfold a protein? Equation 4 contains an expression of electrostatic contribution to the free-energy of folding in terms of $Q^2$ and $Q^3$, the square of the charge on the native and denatured protein, respectively.1051 These terms capture the principle that it is unfavorable to bring a protein’s net charge from the larger volume of the unfolded state to the smaller confines of the native state79,80 (see Figure 7a). This model has been demonstrated to predict the following: (i) the experimentally measured pH–salt phase diagrams for the unfolding of myoglobin, lysozyme, and RNase A14 and (ii) the experimental dependence of the folding free-energy on the square of the net charge.79–82 Equation 4 shows that changing a protein’s charge from $Q$ to $Q \pm 1$, for example from a single oxidative damage event, will change an average protein’s folding stability by $\Delta G(Q) = \Delta G(Q \pm 1) - \Delta G(Q)$, where
Equation 7 is in quantitative agreement with charge-perturbation experiments.\textsuperscript{81,82} It can be computed using only a protein’s sequence. It predicts a proteome-wide distribution of stability changes that is similar to that observed experimentally in point mutations of charged residues, which are reasonable proxies for oxidation.\textsuperscript{83}

A key conclusion from eq 7 is that the change in folding free-energy, \(\Delta \Delta G\), from a damage event will be proportional to the net charge already on the native protein before the damage event. So, any proteins in the proteome that are highly charged and/or relatively unstable to begin with are in greater danger of being destabilized by a single oxidative damage event; see Figure 7b.

Figure 7b shows an interesting implication of the model.\textsuperscript{72} First, the black curve shows the one standard deviation line for the human proteome. It shows that most human proteins are sufficiently neutral to be safe from unfolding by single charge-modification events. Only a few of the proteins in the proteome have a sufficiently high net charge (of either sign) for the destabilization of their native state to be comparable to the stability of some entire proteins (roughly 2–4 \(kT\); see Figure 2).

Now, notice the data points on Figure 7b. These are 20 human proteins known from the literature to be relevant to aging.\textsuperscript{84} These 20 proteins all lie in the high-risk region, and thus, the model predicts that these proteins can be unfolded by a single oxidative charge-modification event. So, changing a single side-chain charge by a random oxidation event could contribute to how aging cells lose protein stability and function.\textsuperscript{85} Figure 8 compares a typical charge distribution...
Figure 9. (a) Foldon Funnel Model predictions for protein folding rates vs number of secondary structure units \( N_s \), compared to data on 93 small single-domain proteins. The inset shows the funnel landscape for this model. (b) Mechanism for how local structures form first and then assemble toward the native state. Reprinted with permission from ref 95. Copyright 2014 American Chemical Society.

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found on the majority of proteins, which are nearly neutral (Figure 8c) and not at risk of unfolding from random oxidation events, with those of highly charged proteins (Figure 8a,b) at high risk of unfolding from single oxidation events.

Additional observations support this mechanism: high net charge is known to predict disorder-prone, unstable proteins, disorder and low stability increase the chance of becoming oxidatively damaged, protein aggregates of old organisms are enriched in damaged proteins, and in budding yeast and worms, aggregates are known to be enriched in highly charged proteins such as ribosomal and DNA-binding proteins. Interestingly, low net charge is also a signature of thermophilic proteins, which face greater stability challenges, as discussed earlier.

DYNAMICAL PROPERTIES OF CELLS ARISE FROM THE FOLDING, SYNTHESIS, DEGRADATION, AND TRANSPORT RATES OF PROTEINS

Below, we review some of the time scales and dynamical processes of proteomes that are important to rapidly duplicating cells.

Protein Folding Happens Fast Enough To Escape the “Grim Reaper” of Proteome Degradation. First, consider the distribution of protein folding times. Experiments show that single-domain proteins fold in vitro over time scales that range over about 8 log orders. Thirumalai developed an early model, predicting that folding rates would scale as \( k_f = k_0 \exp(-N^{1/2}) \) with chain length \( N \). It was remarkably prescient, given the almost complete absence of data at that time. It successfully describes folding rates of proteins and RNA molecules. Recently, a microscopic folding mechanism has been proposed, called the Foldon Funnel Model; see Figure 9.

The model asserts a simple folding mechanism, namely, that local structures form first and rapidly, followed by larger nonlocal structures that assemble more slowly because they have to wait for smaller pieces to form first. The model gives good predictions of folding rates for 93 single-domain proteins from sensible values of helix–coil and hydrophobic interaction parameters (Figure 9a). The model predicts a median nonabundance-weighted folding time of 5 s for the E. coli proteome.

Another model of folding rates is the Topology Polymer Model. It treats the chain conformations more explicitly than the Foldon Funnel Model, fully accounting for entropic costs of chain topological restrictions (see polymer diagrams in ref 94 for details). The Topology Polymer Model also differs by (i) using structure-based domain assignments to predict folding rates and (ii) weighting the folding rates by protein abundance when predicting the proteome folding rate distribution. The Topology Polymer Model gives good predictions for the dependence of folding speed on native topology and unifies different models of folding kinetics. It predicts an average abundance-weighted folding time of 100 ms for the E. coli proteome, and it predicts an average of 170 ms for the yeast proteome. The role of topological constraints in nucleic acids, proteins, and folding kinetics has also been recently revisited using simple folding models. A question for the future remains: What are the folding rates of large single-domain or multidomain proteins? There are not yet many experiments for those types of proteins.

Figure 10a compares the protein folding times for the yeast proteome (from the Topology Polymer Model) with other key rates in the cell. The rate distribution is broad. The most remarkable prediction is that folding speeds seem nearly optimal for outrunning the “grim reaper” of protein degradation, with the slowest-folding proteins just barely out-pacing the fastest protein degradation. This case is made by the black curve in Figure 10a, which is the result of an evolutionary diffusion-drift model of folding rates, resembling the diffusion-drift model of protein stabilities described earlier. The model is based on asserting two physical principles of evolution, namely, that (i) no protein can fold faster than known ultrafast folders, due to conformational speed limits, and (ii) no protein should fold more slowly than the fastest degradation time. Within this interval, the only selection pressure on folding kinetics is simply to “beat the clock” against degradation. When fitted with only one parameter against the folding time distribution derived from the Topology Polymer Model, the model predicts the slowest folding time to be around 10 s. This provides a cushion of an order of magnitude in time separation relative to the fastest degradation times (a few minutes). So, even a protein that degrades at the fastest rate, if not folded off the ribosome by cotranslational folding, has at least a 90% chance of folding before being degraded. For yeast, almost 99% of the proteome’s proteins fold faster than the degradation time (see Figure 10b and ref 96 for details). Among the four outliers, the only protein that folds significantly more slowly has 18 chaperone interaction partners, indicating the important role of chaperones in helping slow folders.

Speed of Cell Duplication Is Limited by the Rate of Protein Translation. What is the speed limit for cell
Protein Translation Speeds Are Limited by Diffusion and Binding. So, why can an amino acid not be added to a growing peptide chain in less than 50 ms in *E. coli*? Translation is known to require several actions: (i) tRNA needs to diffuse to the ribosomal binding site; (ii) the tRNA must settle and bind in the appropriate orientation at this site, with proofreading to verify that it is the correct tRNA; (iii) the peptide is chemically elongated. It is thought that the peptide elongation reaction (iii) is faster than the accommodation step, but this is still debated.115 The rate of tRNA accommodation (ii) has been found experimentally to occur on the same time scale as translation (i) and thus could account for a non-negligible fraction of the total 50 ms. The translation step (i) depends on tRNA concentration. Evidence for its role in a diffusion bottleneck is that cellular tRNA concentrations are roughly the same as those needed to saturate ribosomal kinetics.121 Furthermore, *E. coli* devotes a significant fraction of its dry weight to tRNA (up to 2%121) that could have been spent on more ribosomes, suggesting tRNA plays an important role in protein synthesis speed. Consistent with this, a tRNA diffusion model correctly accounts for the abundance of tRNA with growth rate.121 In short, it appears that the physical processes of translation and the binding and proofreading (ii) are limits to the speed of ribosomal translation.

**Cellular Actions May Be Broadly Rate-Limited by Protein Motions.** Of course, there are very many metabolic rates in the cell. Figure 11a summarizes a broad range of enzyme actions, indicating a predominant time scale around 10−1000 ms.122 What limits their rates? Typical enzyme reactions are often parsed into the following steps:

\[
\text{open} \rightarrow \text{bind} \rightarrow \text{chemical reaction step} \rightarrow \text{close}
\]  

Among these steps, the chemical reaction step itself is often fast. The rate of collision between proteins and small diffusing ligands is on the order of \(10^8 \text{M}^{-1} \text{s}^{-1}\), implying a time scale of 0.1 ms for typical ligand concentrations of 0.1 mM.122 Hence, the rate-limiting steps for enzyme actions appear to be the other steps in eq 9; namely, the opening and closing, binding, product release steps.123−127 These steps can be limited by protein dynamics. Evidence for this view comes from the close correspondence between catalytic rates and the rates of functional motions observed across many proteins, as shown in Figure 11b. However, enzymatic efficiency can be enhanced by other subtle mechanisms as well. For example, binding of allosteric effectors can induce fluctuations128 and alter conformational landscape either by facilitating conformational

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**Figure 10.** (a) Abundance-weighted folding time \((t \text{ in seconds})\) distribution across the yeast proteome (blue) using the topology polymer model,64 which is in good agreement with diffusion-drift model (black) with flat fitness landscape.60 Experimentally measured half-life distribution of the yeast proteome (green)104 shows folding kinetics is faster than protein degradation.60 Median synthesis time is shown in red. (b) The distribution of the ratio of protein half-life and protein folding time.60 Adapted with permission from ref 96. Copyright 2014 Zou et al.

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duplication? In rapidly growing *E. coli* bacteria, DNA replication takes 1−2 ms/base,107 RNA polymerase 10−40 ms/base,108,109 and the ribosome 50 ms/amino acid.110 The ribosome’s slower rate of elongation, combined with its enormous size (since the ribosome itself needs to get copied) and the 10-fold greater cellular abundance of polymerized amino acids relative to nucleotides, makes protein translation the largest bottleneck to cellular growth. In fast-growing *E. coli*, about a third of the cell’s dry weight is ribosome (including rRNA).111,112 What is the maximum rate of protein synthesis? First, cell duplication requires that each ribosome must make a copy of its own proteins. The fastest that a ribosome can copy itself is 6 min, assuming a ribosome’s 7336 amino acids113 are translated at a rate of 20 per second.114 Second, each ribosome must duplicate a corresponding complement of other proteins too. At fast growth rates, an *E. coli* ribosome must make roughly three times its own mass of nonribosomal proteins.111,115 These nearly 30 000 amino acids must be duplicated in series, one-

amino-acid-at-a-time, by each ribosome, predicting a minimum doubling time of 24 min, which approximately equals the observed maximum rate in *E. coli*.111

Interestingly, this 1:3 ratio of ribosomal to nonribosomal proteins also appears to hold in budding yeast, a fast-growing eukaryote.116 So, the minimum cell division time \(t_d\) can be estimated as

\[
t_d = 4rL
\]

where \(r\) is the rate that one ribosome adds one amino acid to a growing protein chain, and \(L\) is the number of amino acids in a ribosome. A ribosome of budding yeast contains 1.6-fold more amino acids than *E. coli’s*113,117 and elongates proteins at half the latter’s speed.110,116 So, if protein translation is indeed the limiting factor in the rate of cell duplication, it implies a minimum doubling time of \(2 \times 1.6 \times 24 \text{ min} = 77 \text{ min}\. This is close to experimental values.118

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transition or altering the width of the free-energy basin and site-specific local flexibility. Partitioning of flux between different pathways can also enhance turnover rates. In spite of these subtleties, the overall role of protein dynamics in enzymatic turnover is clear (Figure 11b). Furthermore, enzyme actions often slow down with increased solvent viscosity (Figure 11c). This is consistent with the observed effect of solvent viscosity on loop closure, which is rate-limiting for catalysis in some enzymes.

So, if cell duplication speeds are ultimately limited by protein motions, why can proteins not wiggle any faster than they do? First, protein conformational energy landscapes are naturally rugged, even along directions of large-amplitude motions. Second, large motions require moving against friction (“wet” friction of the solvent and “dry” friction from internal motions). Third, some motions require local unfolding of secondary structures and that depends on protein folding stability, which is usually marginal. Fourth, the protein conformation that binds the substrate is often little populated, and requires waiting for the right fluctuation. Lastly, there are trade-offs between high affinity for the substrate and stabilization of the transition state conformation.

We hypothesize that biomolecular crowding has two opposing effects on reactions: (i) it increases the concentration of interacting species, but (ii) it hinders and slows the diffusion rate of the reactants. The combination of these two effects predicts a protein diffusion rate $r_d$ that is proportional to $\phi D(\phi)$, where $\phi$ is the protein volume fraction and $D(\phi)$ is the diffusion constant depending on the crowding fraction. The
reduction of diffusion due to volume-excluding monodisperse hard-sphere crowders can be approximated by a simple formula: \( D(\phi) \sim D_0 (1 - \phi / \phi_c)^3 \), where \( D_0 \) is the diffusion in the limit of no crowding, and \( \phi_c \) denotes the volume fraction at which diffusion critically slows down and is estimated to be \( \phi_c \approx 0.58 \).\(^{11,145,146}\) The protein–protein collision rate is

\[
\tau_d \sim \phi (1 - \phi / \phi_c)^3 \tag{10}
\]

Maximizing \( \tau_d \) with respect to \( \phi \) yields the optimal volume fraction of \( \phi_{opt} \approx \phi_c / 3 \approx 0.19 \), close to the typical protein volume fraction (around 0.2) inside a cell.\(^{11,148}\) We can compare this model’s predictions to experiments on bacterial growth rate as a function of salt and crowding volume fraction.\(^{147}\)

To account for heterogeneous protein sizes, two ingredients are needed. First, we have used the hard-particle theory of Minton,\(^{148}\) and its parameters, to estimate how \( D(\phi) \) varies with protein size. This model correctly captures the observed decrease in diffusion with increasing particle size.\(^{148}\) Second, we need to know which particular protein or proteins are responsible for the diffusion limit to cell growth.

Figure 12a shows two different assumptions regarding which proteins are rate-limiting. First, the red curve supposes that all the proteins in the proteome participate in growth, taken by averaging the reaction flux over the molecular weight distribution of the whole \( E. coli \) proteome. Second, an argument has been made\(^{143}\) that one particular type of biomolecule may have an outsized influence on cell dynamics, namely, the tRNA-EF-Tu complex, which are the 70 kDa particles that bring the tRNA molecules to the ribosome in order to elongate the growing peptide chain. As we have argued in the previous section, protein translation, which depends on the rates of amino acid incorporation, may be rate-limiting for cell growth. The basic translation speed of incorporating one amino acid at a time can be further slowed in the presence of crowding due to compromised diffusion. Might the diffusion of the tRNA-EF-Tu complex be growth-limiting? This is a large complex. It will diffuse slowly to the ribosome in the crowded cell environment. This diffusion-bottleneck hypothesis is supported by a recent study showing that ribosomes and tRNA are maintained close to the ratios predicted from diffusion arguments to optimize cell-wide translation rates.\(^{143}\) The black curve in Figure 12 shows the model prediction when the diffusion of tRNA-EF-Tu complexes is considered to be rate-limiting.

Of course, other factors will matter too in the balance of salt and volumes of the cell, including ion fluxes, their regulation, and the balance of ATP.\(^{150}\) The model described above only aims to give a simple estimate of the protein diffusional factor. Cellular crowding is known to affect many physiological processes.\(^{151}\) Crowding can also affect gene expression levels (Figure 12b), reaching a maximum before decreasing at higher densities.\(^{149,152}\) Recent work has also shown cytoplasm can exhibit glassy properties.\(^{153,154}\) The nature of the cytoplasmic environment depends on the size of the cellular objects; for example, small objects experience cytoplasm as a liquid-background while large macromolecules experience a solid-like environment.\(^{154}\) Interestingly, metabolism can also tune the fluidity of the cytoplasm allowing transport of large cellular components that will otherwise be severely constrained in their mobility. Thus, switching between different metabolic states under varying environmental conditions can alter dynamics, cell physiology, and ultimately cellular fitness.\(^{154}\) Figure 12c shows how such relationships represent single slices through a high-dimensional fitness landscape that we are only beginning to understand.

## SUMMARY

While many behaviors of cells emerge from their unique biology, they are fundamentally constrained by the common physics that unites them. Here, we review simple arguments about how these fundamental limits are encoded within the collective physical properties of proteins and proteomes. We describe the role of proteome physics in cell growth laws, providing mechanisms for how cell growth speeds up with temperature and how high salt concentrations slow it down. Electrostatics models give mechanistic insight into the stability gain in thermophiles and the oxidative stability loss in aging and disease. Furthermore, kinetic models of protein folding applied on a global scale show how folding times may be limited by the rate of degradation. And, we note that cell growth appears to be rate-limited by the ribosomal action of adding amino acids to growing protein chains, and by protein motions responsible for enzyme actions. In short, physics can give qualitative and quantitative insights into the growth properties of cells through the use of simple physical models. We believe such global scale models, guided by physicochemical principles, will be increasingly sought after to understand cellular phenotypes and evolution.

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