IGR J16393–4643: a new heavily-obscured X-ray pulsar

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Abstract. An analysis of the high-energy emission from IGR J16393–4643 (= AX J1639.0–4642) is presented using data from INTEGRAL and XMM-Newton. The source is persistent in the 20–40 keV band at an average flux of $5.1 \times 10^{-11}$ ergs cm$^{-2}$ s$^{-1}$, with variations in intensity by at least an order of magnitude. A pulse period of 912.0±0.1 s was discovered in the ISGRI and EPIC light curves. The source spectrum is a strongly-absorbed ($N_H = (2.5 \pm 0.2) \times 10^{23}$ cm$^{-2}$) power law that features a high-energy cutoff above 10 keV. Two iron emission lines at 6.4 and 7.1 keV, an iron absorption edge $\geq 7.1$ keV, and a soft excess emission of $7 \times 10^{-13}$ ergs cm$^{-2}$ s$^{-1}$ between 0.5–2 keV, are detected in the EPIC spectrum. The shape of the spectrum does not change with the pulse. Its persistence, pulsation, and spectrum place IGR J16393–4643 among the class of heavily-absorbed HMXBs. The improved position from EPIC is R.A. (J2000) $= 16^h39^m05.4^s$ and Dec. $= -46^\circ42'12''$ (4'' uncertainty) which is compatible with that of 2MASS J16390535=4642137.

Keywords. Gamma-rays: observations – X-rays: binaries – pulsars: individuals: IGR J16393–4643

1. Introduction

The INTEGRAL core program (CP: Winkler et al. 2003) routinely devotes observation time to Galactic Plane Scans (GPS) and Galactic Centre Deep Exposures (GCDE). These numerous snapshots of the Milky Way can be assembled into mosaic images of long exposure time (~1 Ms). This gamma-ray view of the galaxy, as collected by ISGRI (Lebrun et al. 2003), enabled Bird et al. (2004) to detect 123 sources at a significance above 6$\sigma$. Around 20 of these sources are of unknown origin. A good portion of these new sources probably belong to the class of heavily-absorbed High Mass X-ray Binaries (HMXBs) that are concentrated along the galactic plane and in the spiral arms.

High-Mass X-ray Binaries are composed of a compact object such as a neutron star or a black hole that orbits a massive stellar companion. Depending on the type of companion, known HMXBs can be divided into two groups (van Paradijs 1983). Most HMXBs classified by Liu et al. (2000) contain a Be star. These systems are usually transient sources with hard spectra. The compact object has a wide orbit which mostly keeps it away from the Be star and its disk. Outbursts in these systems are due primarily to the compact object approaching the star and accreting matter from the slow, dense stellar wind. The second group of HMXBs features an O or B supergiant star. The orbit of the compact object places it well within the stellar wind, so material from the supergiant can be fed directly to the compact object through Bondi accretion, or it can pass to the compact object via an accretion disk. The latter mechanism is typically found in bright X-ray binaries in which the Roche lobe overflow of gas from the OB star supercedes the flow of accreting matter. For less luminous binaries, the OB star does not fill its Roche lobe and the behaviour of the X-ray source is determined predominantly by the stellar wind. X-ray emission in supergiant HMXB systems is usually persistent, with flares stemming from inhomogeneities in the wind. Neutron stars with strong magnetic fields develop a hot spot for accretion which can result in a pulsation.

The number of persistent and heavily-absorbed HMXBs associated with supergiant companions has increased thanks to deep, wide-field observations by INTEGRAL combined with follow-up X-ray monitoring by XMM-Newton (Walter et al. 2003, Rodriguez et al. 2003, Hill et al. 2005, Lutovinov et al. 2005b). So far, these sources have been detected preferentially in the Norma Arm region (Walter et al. 2004) which features high formation rates of OB supergiant stars. It is there that INTEGRAL uncovered its first new source,
IGR J16318−4848 (Walter et al. 2003). The Norma Arm also harbors IGR J16393−4643. This object was initially discovered in the X-ray band and listed as AX J163904−4642 in the ASCA Faint Source Catalog (Sugizaki et al. 2001). Its flat power-law slope, its absorption ($N_H = 13^{+15}_{-7} \times 10^{22} \text{ cm}^{-2}$), its lack of radio emission, and its position in the galactic plane compelled the authors to classify it as a HMXB. A non-thermal radio counterpart was recently detected in the 4702 (Hartman et al. 1999) noted by Malizia et al. 4848 (Walter et al. 2003). The Norma Arm also harbors IGR J16393−4643 in the Faint Source Catalog (Sugizaki et al. 2001). Its flat power-law slope, its absorption ($N_H = 13^{+15}_{-7} \times 10^{22} \text{ cm}^{-2}$), its lack of radio emission, and its position in the galactic plane suggested a heavily-obscured microquasar interpretation for the HMXB (Combi et al. 2004), and would help explain which suggests a dust-enshrouded microquasar interpretation for the unidentified EGRET source 3EG J1639−4702 (Hartman et al. 1999) noted by Malizia et al. (2004).

We detected IGR J16393−4643 in the ISGRI GPS data of the Norma Arm’s first visibility period, and we obtained a follow-up observation with XMM-Newton. The set of INTEGRAL and XMM-Newton data are introduced in Section 2.2. In Section 2.3 the refined X-ray position is used to locate the most likely counterpart at other wavelengths. Timing and spectral analyses are presented in Sections 3.1 and 3.2 respectively. Finally, we discuss the nature of the source and we offer our conclusions in Section 4.

2. Observations and Data Sets

2.1. INTEGRAL Data and Imaging

The INTEGRAL data consist of all CP pointings during revolutions 30–260, as well as pointings that were public by January 3, 2005, which had the source within the ISGRI field of view (FOV). To improve the quality of the output mosaic, pointings with exposure times below 1 ks were ignored. The resulting data set groups roughly 1500 pointings with an average exposure time of 2 ks each.

Version 4.2 of the INTEGRAL Offline Scientific Analysis (OSA) software was used to reduce raw data into images. Source extraction employed version 18 of the INTEGRAL General Reference Catalog (Ebisawa et al. 2003) selected for objects detected by ISGRI. These tools are available to the public through the INTEGRAL Science Data Centre (ISDC: Courvoisier et al. 2003).

Intensity, significance, variance, and exposure mosaic images were constructed from background-subtracted images of individual pointings. Figure 1 provides an example of an intensity map of IGR J16393−4643 and its vicinity in the 20–40 keV band from the pointings of revolutions 30 to 185. The effective exposure time is 670 ks after correcting for vignetting. Using this image, we extracted a source location of R.A. (J2000)=$16^339^005^a$ and Dec.=$-46^142^3$ (26” uncertainty) which agrees with and improves the ISGRI position of Bird et al. (2004). The mean flux (20–60 keV) of the source is $0.73 \pm 0.02$ counts per second (cps), or 4.9 mCrab, at a significance of 36x.

2.2. XMM-Newton Data and Imaging

XMM-Newton (Jansen et al. 2001; Strüder et al. 2001; Turner et al. 2001) observed IGR J16393−4643 on March 21, 2004, from 08:21:15 to 11:41:15 (UT). We used the Science Analysis System (SAS) v. 6.1.0 software to analyse the data and to extract the EPIC spectrum. The data were screened for background variability by excluding time intervals in which the count rate above 10 keV was greater than a selected threshold (0.33 cps for EPIC/MOS1, 0.25 cps for EPIC/MOS2, and 2.4 cps for EPIC/PN). After screening, the effective exposure times were around 8.4, 8, and 7 ks for EPIC/MOS1, EPIC/MOS2, and EPIC/PN, respectively.

IGR J16393−4643 is clearly detected in images taken by EPIC. Figure 2 presents an image of IGR J16393−4643 from the EPIC/MOS1 camera. Furthermore, the source coincides with the updated ISGRI position and error circle from this paper. The refined X-ray position averaged from EPIC/MOS and EPIC/PN is R.A. (J2000)=$16^339^005.4^a$ and Dec.=$-46^142^12^"$ (4” uncertainty).
We searched for counterparts at other wavelengths and found 3. Counterparts
the refined ISGRI error circle from this paper is superimposed.

3. Counterparts

We searched for counterparts at other wavelengths and found a single infrared source belonging to the Two Microns All-Sky Survey (2MASS: Cutri et al. 2003) in the XMM-Newton error box. This potential counterpart, 2MASS J16390535–462137, is located about 2' away from the XMM-Newton position at R.A. (J2000) = 16h39m05.36s and Dec. = −46°42′13.7″ (0.06″ uncertainty). It appears in the J, H, and Ks bands with magnitudes of 14.63±0.06, 13.02±0.04, and 12.78±0.04, respectively

A possible association of IGR J16393−4643 with the EGRET source 3EG J1639−4702 (Hartman et al. 1999) has been proposed by Malizia et al. (2004). Another ISGRI source, IGR J16358−4726, and 4U 1630−47 are just outside the EGRET 95% error contour. In the 2° × 5° degree section of the Norma Arm presented in Fig. 1, the probability to observe a counterpart in the 0.56″ error radius of the EGRET source is close to 1, so the coincidence with IGR J16393−4643 is probably a chance one.

4. Timing Analysis

4.1. Long-term Variability

Most ISGRI pointings have exposure times that are too brief for a significant detection. Therefore, a mosaic was prepared for each 3-day spacecraft revolution (from 30–244) in the 20–60 and 60–150 keV bands. We extracted the source flux, error, and significance by fitting a Gaussian with a fixed point spread function to the mosaic images. Upper limits (3σ), calculated from the variance maps, are provided when the source is not detected at the 3σ level in a mosaic image with an effective exposure time above 7 ks. There were no detections in mosaic images of the 60–150 keV band so its light curve is omitted.

Figure 3 illustrates the persistence of IGR J16393−4643 in the 20–60 keV energy range. Table I collects the 15 revolutions in which the source is detected in the mosaic image with a significance above 4σ. The average flux in these mosaics is 0.86 cps (5.6 mCrab) with a cumulative exposure time of 1.5 Ms. The mean flux per revolution varies by a factor of at least 6 from 0.21±0.05 counts per second (cps) or 1.4 mCrab in revolution 100 (205 ks), to 1.39±0.14 cps (9.3 mCrab) in revolution 65 (196 ks). When the flux value during revolution 100 is compared to the highest flux registered in a single 2-ks pointing of the 20–60 keV band (4.3±0.6 cps, or 29 mCrab, during MJD 52673.623–52673.641), we find that the source flux varies by a factor larger than 20.

4.2. Pulsations

During the EPIC/PN observation of 8 ks, the source count rate varied from 1 to 4 cps (Fig 5). Furthermore, the variations occur periodically. By searching for modulations in the power spectrum and in the χ2 distribution, we obtain a period of 912±5 s at a χ2 of ~600 for 30 bins per period (Fig 6b). The pulse period is determined by the centre of a Gaussian fit to the χ2 distribution, and the error is calculated using the procedure developed by Horne & Baliunas (1986) on Lomb-Scargle periodograms. The EPIC/PN pulse profile folded with a period of 912 s, illustrated in Fig 7b, indicates a pulse fraction of 38±5%.
For ISGRI, the pulse is best detected in the 100-s light curve of revolution 38 in the 15–40 keV band since this revolution has enough effective exposure time (50 ks) with low off-axis angles (~ 4°), and its light curve is free of gaps that hinder a periodicity search. Figure 5 presents the $\chi^2$ distribution centered at 912 s, with 9 bins per period. The phase diagram is shown in Fig. 6b, and the pulse fraction is $57 \pm 24\%$. We merged the 100-s light curves (15–30 keV) from revolutions 37–55, and derived a pulse period and error of $912.0 \pm 0.1$ s.

With respect to the pulse fractions, the amplitude does not change significantly as a function of energy. Both display a jagged rise to a peak flux followed by a drop. There appear to be primary and secondary peaks in the pulse profiles of both instruments at phases of $\sim 0.7$ and $\sim 0.2$, respectively. Although the ISGRI and EPIC periods are derived from observations about 400 days apart, the period is not accurate enough to search for possible variations. In addition, the ISGRI light curves of revolutions that are nearly concurrent with the EPIC observation present low source significances which makes it difficult to search for a period.

There is an indication in the EPIC light curve (Fig. 5) that the source varies over timescales longer than the pulse period. At the beginning of the observation, the maximum of the pulse is about 3 cps. This maximum rises to 4 cps about 5 ks after the start time, and then decreases.

We did not detect an orbital period of the order of a few days in the combined ISGRI light curves (15–30 keV, 100-s resolution) of detected revolutions between 37 and 55.

### 5. Spectral Analysis

#### 5.1. Average Spectrum

To extract an ISGRI spectrum, we created mosaics accumulating the data from revolutions 37–185 in the following energy bands: 22.09–25.92, 25.92–30.23, 30.23–40.28, 40.28–51.29, and 51.29–63.26 keV. These boundaries were chosen to conform with the current response matrices. The source is clearly detected up to 40 keV, so spectral bins at higher energies are ignored. A power law fit to the ISGRI spectrum returns a relatively steep photon index of $4.5 \pm 0.4 \ (\chi^2 = 0.32)$.

Spectral extraction for EPIC/PN, EPIC/MOS1, and EPIC/MOS2 relied on single and double events within a cir-
Fig. 6. Top (a): Period search ($\chi^2$ distribution) on the ISGRI light curve (15–40 keV) of revolution 38, centred at 912 s (vertical line), with 9 bins per period, and a resolution of 0.4 s. Bottom (b): Pulse profile of the folded ISGRI light curve (15–40 keV) for a period of 912 s. The start time is MJD 52674.53833.

Fig. 7. Top (a): Distribution of $\chi^2$ for the EPIC/PN 2–10 keV light curve with 30 bins per period, a resolution of 1 s, and centred at 912 s (vertical line). Bottom (b): Folded EPIC/PN light curve for a pulse period of 912 s beginning at MJD 53085.38738. The intervals indicate the pulse and plateau states of the phase-resolved spectrum.

circle of radius 25′′ centred on the source. To estimate the background, we made an event list from a circle of equal radius in the same detector and at an equivalent distance from the detector’s edge. Channels were configured to collect 30 counts per bin. An absorbed power law fit to the EPIC spectra has a photon index of 1.0±0.1 with a column density of 3.1±0.1×10^{23} cm^{-2} ($\chi^2_\nu = 1.15$).

The EPIC spectrum is heavily-absorbed below 4 keV. Iron emission lines appear at 6.4 keV (Fe Kα) and at 7.1 keV (Fe Kβ). A discontinuity in the continuum above 7.1 keV suggests an iron absorption edge. There is an indication of an excess of soft emission between 0.5 and 2 keV. Including a soft blackbody component in an absorbed power law adjusted to the EPIC spectra decreases the $\chi^2$ by 9 for 680 degrees of freedom (dof), or a 0.5% probability that this feature is due to chance. Using a partial covering absorption instead of a blackbody raises the $\chi^2$ by 60 for the same dof (679). A partially-ionised absorber is also insufficient to model the soft excess.

Simply fitting an absorbed power law to the spectra from EPIC and ISGRI yields $N_H \sim 4 \times 10^{23} cm^{-2}$ and $\Gamma \sim 2.3$, but the fit is poor ($\chi^2_\nu = 2.1$) with significant residuals in the hard X-rays. The EPIC photon index of 1.0±0.1 is smaller than the one observed for ISGRI (4.5±0.4) which indicates a spectral break. Such spectral shapes—a flat power law at low energies with a high-energy cutoff between 10–20 keV beyond which the slope steepens—are typical of X-ray pulsars [White et al. 1995].
We fit the combined EPIC and ISGRI spectra (Fig. 3) with a broken power law (BP), a power law with an exponential cutoff energy (CP), and a Compton emission (CE) model (\texttt{comptt} within Xspec). The models are modified by a galactic absorption in the direction of the source (fixed at $2.2 \times 10^{22}$cm$^{-2}$ (Dickey & Lockman, 1990), and an intrinsic photoelectric absorption with free iron abundances ($Z_{Fe}$). Soft excess emission is represented by a blackbody (with a fixed temperature $kT_B = 0.06$ keV) that is affected by the galactic absorption only. Two narrow Gaussians (with widths fixed at 0) describe the iron lines. A constant ($C_I$) accounts for the asynchronous observations and for cross-calibration uncertainties.

The BP model ($\chi^2_{red}$/dof = 0.95/675) has photon indices of $\\Gamma_1 = 0.9^{+0.1}_{-0.2}$ and $\\Gamma_2 = 4.6^{+0.8}_{-0.3}$. However, the BP model does not constrain the break energy ($E_{\text{break}} > 17$ keV) nor $C_I$ because it requires one more parameter than the other models. The CP model ($\chi^2_{red}$/dof = 1.07/676) gives $\\Gamma = 0.8 \pm 0.2$, and a cutoff energy of $E_{\text{cut}} = 10 \pm 1$ keV, but residuals remain at high energy because of the low cutoff temperature. In Lutovinov et al. (2005), an absorbed cutoff power law fit to the combined ASCA and ISGRI spectra had a comparable cutoff energy of $11 \pm 1$ keV, but the photon index was poorly constrained ($\\Gamma = 1.3 \pm 1.0$).

Parameters from the CE model ($\chi^2_{red}$/dof = 0.95/676) are listed in Table 2. The column density ($N_H$) is $(25 \pm 2) \times 10^{22}$cm$^{-2}$. The comptonising medium has an electron temperature ($kT_e$) of $4.4 \pm 0.3$ keV with an optical depth ($\tau$) of 9.6. The Fe K$_\alpha$ line is at 6.41$^{±0.05}$ with an optical depth ($\tau$) of 0.03 and has an equivalent width ($W$) of 6.41$^{±0.05}$ keV. The detection of the K$_\beta$ line is marginal given the uncertainties and its F-test probability of 3%. However, the ratio of iron intensities ($F_{K\beta}/F_{K\alpha}$) is consistent with the value expected from the photoionisation of iron (Kaastra & Mewe, 1993). The observed soft excess flux (in units of $10^{-13}$ergs cm$^{-2}$ s$^{-1}$) is 4.4 in the 2–10 keV band, and 5.1 in the 20–60 keV band. The observed soft excess flux between 0.5 and 2 keV is $7 \times 10^{-13}$ergs cm$^{-2}$ s$^{-1}$. Since $C_I \sim 1$, the EPIC observation occurred during a period in which the source was in an average state.

5.2. Phase-resolved Spectrum

The EPIC/PN spectra were binned according to two states: a pulse state around the primary peak in the folded light curve (phase 0.43–0.77 in Fig. 7b), and a plateau state (phases 0.00–0.07 and 0.87–1.00). Average count rates for the pulse and plateau states are almost a factor of 2 apart at $\sim1.2$ and $\sim0.7$, respectively. The CE model was applied to the phase-resolved EPIC spectra constrained with the ISGRI spectrum. Bins for the plateau gather 20 counts and bins for the pulse collect 40 counts per channel. The soft excess is not prominent in either spectrum so the blackbody component is omitted. This prevents an evaluation of the influence of the pulse on the soft excess. Table 2 reports parameters for the CE model fit to the phase-resolved spectra. The pulsation affects the normalisations but does not modify the shape of the spectra nor its parameters, specifically the $N_H$, in any appreciable way.

6. Discussion

Observations of IGR J16393–4643 by INTEGRAL and XMM-Newton present a source that is highly-obscured and persistent, with an intensity varying by a factor larger than 20. The source has a pulsation period of 912.0±0.1 s. These attributes, along with its spectral shape, its neutral iron lines, and its lack of radio emission, suggest a HMXB system consisting of a magnetised neutron star orbiting inside the dense stellar wind of a supergiant companion. In the 2–10 keV energy band, the unabsorbed flux of IGR J16393–4643 is $9.2 \times 10^{-11}$ ergs cm$^{-2}$ s$^{-1}$. Assuming a luminosity of $1.2 \times 10^{36}$ ergs s$^{-1}$, which is typical for accretion-driven X-ray pulsars (Bildsten et al., 1997), the source is located approximately 10 kpc away.

The absorbing column density is large whenever the source is detected, whether by ASCA ($N_H = 13.2^{+2}_{-1} \times 10^{22}$cm$^{-2}$), or now with INTEGRAL and XMM-Newton ($N_H = (25\pm2) \times 10^{22}$cm$^{-2}$). The strong absorption below 5 keV is intrinsic since it is an order of magnitude larger than the galactic absorption along the
line of sight ($N_\text{H} = 2.2 \times 10^{22} \text{cm}^{-2}$). The absorption does not vary with the pulsation, which means that it is not related to the accretion column, but it may change with the orbital phase. This absorption indicates the presence of optically-thick material surrounding the compact object, which is consistent with the detection of iron lines.

Emission lines at 6.4 and 7.1 keV trace the quantity of matter in the shell that envelopes the neutron star. These lines are at the positions that would be expected from the fluorescence of cold neutral matter illuminated by continuum X-rays, and their equivalent widths are compatible with a spherical distribution of dense gas around the compact object (Matt 2002). A Kα line at 6.41 ± 0.03 keV (≥1.925 Å) corresponds to an ionisation of at most Fe XVIII (House 1969), which does not constrain the distance between the ionising source and the inner surface of the ionised shell. Given the errors on the line equivalent widths, we can determine whether the fluxes of the iron lines are modified by the pulse.

There appears to be a soft excess flux of $7 \times 10^{-15} \text{ergs cm}^{-2} \text{s}^{-1}$ that is best represented by a blackbody rather than with a partial covering or partially-ionised absorber. The scattering of X-rays by the stellar wind is the most likely explanation for the soft excess emission (Haberl & White 1990). To what extent the pulse affects the soft excess emission is still unknown as the soft excess flux is not prominent in the phase-resolved spectrum.

The ASCA Faint Source Catalog (Sugizaki et al. 2001) listed IGR J16393–4643 among the brightest objects in its catalog, but the uncertainty of the position and the heavy absorption prevented (Sugizaki et al. 2001) from associating this source with an optical counterpart. Malizia et al. (2004) noted the possible association between IGR J16393–4643 and the unidentified EGRET source 3EGJ1639–4702 (Hartman et al. 1999). Non-thermal radio emission was recently detected within the ASCA error box (Combi et al. 2004) which suggests a dust-enshrouded microquasar and could help justify the IBIS and EGRET association. However, the XMM-Newton position (R.A. (J2000)= $16^h39^m05.4^s$, Dec,= $-46^\circ42'12''$, 4'' uncertainty) is incompatible with the position of the radio source MOST J1639.0–4642, or any known radio source from the Vizier database. The pulsation and the lack of radio emission invalidate the microquasar interpretation for the HMXB. Also, the chance association of the EGRET source with IGR J16393–4643 is strong given the high density of sources in the region.

A single infrared candidate, 2MASS J16390535–4642137, lies within the XMM-Newton error circle of IGR J16393–4643. The high-energy spectral and timing characteristics of IGR J16393–4643 are reminiscent of other heavily-absorbed, wind-accreting X-ray pulsars with OB supergiant companions (Walter et al. 2004). A Be stellar companion is unlikely given that such systems are usually transient. Infrared observations should confirm the supergiant nature of the companion. If this is the case, it will constitute another argument rejecting the EGRET association (Orellana & Romero 2005).

Recently, combined INTEGRAL and XMM-Newton observations of the Norma Arm have revealed more pulsating X-ray binaries than were previously known. These objects have joined the ranks of what might be a new class of heavily-absorbed HMXBs that were previously undetected below 5 keV. The increasing sample size these objects represent should enable meaningful statistical studies to be performed. Understanding the nature of sources such as IGR J16393–4643 could shed light on the structure of stellar winds, provide constraints on the masses of neutron stars, and help elucidate the evolution of binaries.

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Table 2. Parameters from the Compton emission (CE) model fit to the spectra of EPIC and ISGRI combined. Also listed are the values of the CE model (without a blackbody) fit to the pulse and plateau states of the phase-resolved EPIC/PN spectra constrained with the ISGRI spectrum. The integrated fluxes are listed as observed ($F_{2-10}$) or unabsorbed ($F_{2-10}^{\text{UA}}$). Errors denote 90% confidence.

| Parameter | CE       | Pulse   | Plateau  | Unit          |
|-----------|----------|---------|----------|---------------|
| $C_I$     | 0.8 ± 0.2| 0.6 ± 0.2| 1.3 ± 0.7|               |
| $\tau$    | 9 ± 1    | 12±6    | 8±4      |               |
| $kT_e$    | 4.4 ± 0.3| 4.4±0.3 | 4.4 ± 0.4| keV           |
| $N_H$     | 25±2     | 23±2    | 24±3     | $10^{22} \text{cm}^{-2}$ |
| $Z_N$     | $1.7^{+0.4}_{-0.2}$ | $0.9^{+0.6}_{-0.3}$ | $1.3^{+0.8}_{-0.6}$ | $Z_\odot$      |
| $F_{2-10}$| 4.4      | 5.2     | 3.0      | $10^{-11} \text{ergs cm}^{-2} \text{s}^{-1}$ |
| $F_{2-10}^{\text{UA}}$| 9.2      | 9.8     | 6.7      | $10^{-11} \text{ergs cm}^{-2} \text{s}^{-1}$ |
| $E_{K\alpha}$| 6.41±0.03| 6.4±0.03| 6.41±0.04| keV           |
| $F_{K\alpha}$| 1.1±0.2 | 1.1±0.4 | 0.8±0.4  | $10^{-3} \text{ph cm}^{-2} \text{s}^{-1}$ |
| $E_{K\beta}$ | 7.1±0.1 | 7.1±0.3 | 7.2±0.3  | keV           |
| $F_{K\beta}$ | <0.62   | <1.1    | <1.3     | $10^{-3} \text{ph cm}^{-2} \text{s}^{-1}$ |
| $\chi^2/\text{d.o.f.}$ | 0.95/676 | 0.92/140 | 0.929/4    |               |
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