A quantum Goldreich-Levin theorem with cryptographic applications

Mark Adcock∗ Richard Cleve†

Department of Computer Science
University of Calgary
Calgary, Alberta, Canada T2N 1N4

Abstract
We investigate the Goldreich-Levin Theorem in the context of quantum information. This result is a reduction from the computational problem of inverting a one-way function to the problem of predicting a particular bit associated with that function. We show that the quantum version of the reduction—between quantum one-way functions and quantum hard-predicates—is quantitatively more efficient than the known classical version. Roughly speaking, if the one-way function acts on $n$-bit strings then the overhead in the reduction is by a factor of $O(n/\varepsilon^2)$ in the classical case but only by a factor of $O(1/\varepsilon)$ in the quantum case, where $\frac{1}{2} + \varepsilon$ is the probability of predicting the hard-predicate. Moreover, we prove via a lower bound that, in a black-box framework, the classical version of the reduction cannot have overhead less than $\Omega(n/\varepsilon^2)$.

We also show that, using this reduction, a quantum bit commitment scheme that is perfectly binding and computationally concealing can be obtained from any quantum one-way permutation. This complements a recent result by Dumais, Mayers and Salvail, where the bit commitment scheme is perfectly concealing and computationally binding. We also show how to perform qubit commitment by a similar approach.

1 Introduction

Fast quantum algorithms are potentially useful in that, if quantum computers that can run them are built, they can then be used to solve computational problems quickly. Algorithms can also be the basis of reductions between computational problems in instances where the underlying goals are different from fast computations. For example, reductions are often used as indicators that certain problems are computationally hard, as in the theory of NP-completeness (see [12] and references therein). Another domain where reductions play an important role is in complexity-based cryptography, where a reduction can show that breaking a particular cryptosystem is as difficult (or almost as difficult) as solving a computational problem that is presumed to be hard.

We investigate such a cryptographic setting where quantum algorithms yield different reductions than are possible in the classical case: the so-called Goldreich-Levin Theorem [13]. This result is a reduction from the computational problem of inverting a one-way function to the problem of predicting a particular hard-predicate associated with that function. Roughly speaking, a one-way function is a function that can be efficiently computed in the forward direction but is hard to

∗Email: mark.adcock@ucalgry.com
†Email: cleve@cpsc.ucalgary.ca. Partially supported by Canada’s NSERC.
compute in the reverse direction, and a hard-predicate of a function is a bit that can be efficiently computed from the input to the function and yet is hard to estimate from the output of the function. We show that the quantum version of the reduction is quantitatively more efficient than the known classical version. Moreover, we prove via a lower bound that, in a black-box framework, the classical version of the reduction cannot be made as efficient as the quantum version.

Goldreich and Levin essentially showed that, for a problem instance of size $n$ bits, if their hard-predicate can be predicted with probability $\frac{1}{2} + \epsilon$ with computational cost $T$ then the one-way function can be inverted with computational cost $O(T D(n,\varepsilon))$, where $D(n,\varepsilon)$ is polynomial in $n/\varepsilon$. Taken in its contrapositive form, this means that, if inverting the one-way function requires a computational cost of $\Omega(T)$, then predicting the hard-predicate with probability $\frac{1}{2} + \epsilon$ requires a computational cost of $\Omega(T/D(n,\varepsilon))$. Note that if we start with a specific lower bound of $\Omega(T)$ for inverting the function then we end up with a weaker lower bound—by a dilution factor of $D(n,\varepsilon)$—for breaking the hard-predicate. In [14], it is shown that the dilution factor can be as small as $O(n/\varepsilon^2)$.

We show that there is a quantum implementation of the reduction where the dilution factor is only $O(1/\varepsilon)$. We also show that $\Omega(n/\varepsilon^2)$ is a lower bound on the dilution factor for any classical implementation of the reduction in a black-box framework. In the standard parameterization of interest in cryptography, $T$ is assumed to be superpolynomial in $n$ and $\varepsilon \in 1/n^{O(1)}$. In this case, although $1/\varepsilon$ is smaller than $n/\varepsilon^2$, the diluted computational cost, $T/D(n,\varepsilon)$, remains superpolynomial in both cases. However, there are other parameterizations where the difference between the achievable quantum reduction and best possible classical reduction is more pronounced. One example is the case where $T = n^3$ and $\varepsilon = 1/n$. If we start with a classical one-way function that requires a computational cost of $\Omega(n^3)$ to invert and apply the Goldreich-Levin Theorem to construct a classical hard-predicate then the reduction implies only that the computational cost of predicting the predicate with probability $\frac{1}{2} + \frac{1}{n}$ is lower bounded only by a constant. However, if we start with a quantum one-way function that requires a computational cost of $\Omega(n^3)$ to invert and apply our quantum version of the Goldreich-Levin Theorem then the computational cost of predicting the predicate with probability $\frac{1}{2} + \frac{1}{n}$ is lower bounded by $\Omega(n^2)$.

A particular application of hard-predicates is for bit commitment. Recall the now well-known result that an information theoretically secure bit commitment scheme cannot be based on the information-theoretic properties of quantum devices alone [16, 17]. Of course, this is also the case with classical devices, though computationally secure bit commitment schemes have been widely proposed, investigated, and applied. Such schemes can be based on the existence of one-way permutations. Most of these proposed one-way permutations are hard to invert only if problems such as factoring or the discrete logarithm are hard, and are insecure against quantum computers, which can efficiently solve such problems [18]. Recently, Dumais, Mayers and Salvail considered the possibility of quantum one-way permutations [14], and showed how to base quantum bit commitment on them (see also [10]). Their scheme is perfectly concealing and computationally binding, in the sense that changing a commitment is computationally hard if inverting the permutation is hard. We exhibit a complementary quantum bit commitment scheme that is perfectly binding and computationally concealing. As with hard-predicates, the dilution factor in the measure of computational security is lower than possible with the corresponding classical construction. Furthermore, a possible advantage of our protocol is that the information that must be communicated and stored between the parties consists of $O(n)$ classical bits for bit commitment (and $O(n)$ classical bits plus one qubit for qubit commitment), whereas the scheme in [11] employs $O(n)$ qubits.

The organization of this paper is as follows. In Section 2, we investigate a simple black-box
problem that is related to the Goldreich-Levin Theorem. In Section 3, we give definitions pertaining to one-way permutations and hard-predicates (classical and quantum versions) and investigate the complexity of reductions from the former to the latter (applying results from Section 2). In Section 4, we show how to use the Goldreich-Levin Theorem to construct a perfectly binding and computationally concealing quantum bit commitment scheme from a quantum one-way permutation.

2 A black-box problem

Our results about the Goldreich-Levin Theorem (which are in Section 3) are based on the query complexity of the following black-box problem, which we refer to as the GL problem (see, e.g., [2]). Let $n$ be a positive integer and $\varepsilon > 0$. Let $a \in \{0,1\}^n$ and let information about $a$ be available only from inner product and equivalence queries, which are defined below in the classical case (and later on generalized to the case of quantum information).

Definition 1 A classical inner product (IP) query (with bias $\varepsilon$) has input $x \in \{0,1\}^n$ and outputs a bit that is slightly correlated with $a \cdot x$ (the inner product of $a$ and $x$ modulo two) in the sense that

$$\Pr_{x}[IP(x) = a \cdot x] \geq \frac{1}{2} + \varepsilon.$$ 

The above probability is with respect to a random $x \in \{0,1\}^n$.

Definition 2 An quantum equivalence (EQ) query has input $x \in \{0,1\}^n$, and the output is 1 if $x = a$ and 0 otherwise.

The goal is to determine $a$ with a minimum number of IP and EQ queries. A secondary resource under consideration is the number of auxiliary bit/qubit operations. It should be noted that, when $\varepsilon = \frac{1}{2}$, this is essentially equivalent to a problem that Bernstein and Vazirani considered, where IP queries return $a \cdot x$ on input $x$. For this problem, $n$ IP queries are necessary and sufficient to solve it classically; however, it can be solved with a single (appropriately defined) quantum IP query. (See also [19].) When $\varepsilon$ is small—say, $\varepsilon \in 1/n^{O(1)}$—an efficient classical solution to this problem is nontrivial. The correctness probability of an IP query for a particular $x$ cannot readily be amplified by simple techniques such as repeating queries; for some $x$, $IP(x)$ may always be wrong. Goldreich and Levin [13] were the first to (implicitly) solve this problem with a number of queries and auxiliary operations that is polynomial in $n/\varepsilon$—and this is the basis of their cryptographic reduction in Theorem 4.

We show that any classical algorithm solving the GL problem with constant probability must make $\Omega(n/\varepsilon^2)$ queries (for a reasonable range of values of $\varepsilon$), whereas there is a quantum algorithm that solves the GL problem with $O(1/\varepsilon)$ queries. For the quantum version of the GL problem, quantum IP and EQ queries are defined (in Definitions 3 and 4) as unitary operations that correspond to Definitions 1 and 2 in a natural way. We begin with the classical lower bound.

1Unless otherwise specified, a “random” element of a set means with respect to the uniform distribution.
Theorem 1 Any classical probabilistic algorithm solving the GL problem with success probability \( \delta > 0 \) requires either more than \( 2^{n/2} \) EQ queries or \( \Omega(\delta n/\varepsilon^2) \) IP queries when \( \varepsilon \geq \sqrt{n/2} - n/3 \).

Proof: The proof uses classical information theory, bounding the conditional mutual information about an unknown string that is revealed by each IP query, in conjunction with an analysis of the effect of EQ queries.

It is useful to consider an algorithm to be successful on a particular input if and only if it performs an EQ query whose output is 1 (at which point the value of \( a \) has been determined).

We begin by showing that it is sufficient to consider algorithms (formally, decision trees) that are in a convenient simple form. First, by a basic game-theoretic argument, it suffices to consider deterministic algorithms, where their input data—embodied in the black-boxes for IP and EQ queries—may be generated in a probabilistic manner. Second, it can be assumed that all EQ queries occur only after all IP queries have been completed. To see why this is so, start with an algorithm that interleaves IP and EQ queries, and modify it as follows. Whenever an EQ query occurs before the end of the IP queries, the modified algorithm stores the value of the input to the query and proceeds as if the result were 0. Then, at the end of the IP queries, each such deferred EQ query is applied. The modified algorithm will behave consistently whenever the actual output of a deferred EQ query is 0, and also it will perform (albeit later) any EQ query where the output is 1. Henceforth, we consider only algorithms with the above simplifications.

Now we describe a probabilistic procedure for constructing the black boxes that perform IP and EQ queries. First, \( a \in \{0, 1\}^n \) is chosen randomly according to the uniform distribution. Then a set \( S \subseteq \{0, 1\}^n \) is chosen randomly, uniformly subject to the condition that \( |S| = (\frac{1}{2} + \varepsilon)2^n \) (assuming that \( \varepsilon 2^n \) is an integer). Then

\[
IP(x) = \begin{cases} 
  a \cdot x & \text{if } x \in S \\
  a \cdot \overline{x} & \text{if } x \notin S 
\end{cases}
\]  

and

\[
EQ(x) = \begin{cases} 
  1 & \text{if } x = a \\
  0 & \text{if } x \neq a.
\end{cases}
\]

Consider an algorithm that makes \( m \) IP queries. If \( m \geq \delta n/\varepsilon^2 \) then the theorem is proven. Otherwise, since \( \varepsilon \geq \sqrt{n/2} - n/3 \), we have

\[
m < \frac{\delta n}{\varepsilon^2} \leq \delta 2^{2n/3}.
\]  

We proceed by determining the amount of information about \( a \) that is conveyed by the application of \( m \) IP queries. Let \( A \) be the \( \{0, 1\}^n \)-valued random variable corresponding to the probabilistic choice of \( a \in \{0, 1\}^n \), and let \( Y_1, Y_2, \ldots, Y_m \) be the \( \{0, 1\} \)-valued random variables corresponding to the respective outputs of the \( m \) IP queries. Let \( H \) be the Shannon entropy function (see, e.g., \[4\]). Then, for each \( i \in \{1, 2, \ldots, m\} \),

\[
H(A|Y_1, Y_2, \ldots, Y_i) = H(A|Y_1, \ldots, Y_{i-1}) - H(Y_i|Y_1, \ldots, Y_{i-1}) + H(Y_i|A, Y_1, \ldots, Y_{i-1}).
\]  

(5)

Combining the above equations yields

\[
H(A|Y_1, Y_2, \ldots, Y_m) = H(A) + \sum_{i=1}^{m} (H(Y_i|A, Y_1, \ldots, Y_{i-1}) - H(Y_i|Y_1, \ldots, Y_{i-1})).
\]  

(6)
We shall now bound each term on the right side of Eq. 3. Since the a priori distribution of $A$ is uniform, $H(A) = n$. Also, since the entropy of a single bit is at most 1, $H(Y_i|Y_1, \ldots, Y_{i-1}) \leq 1$ for all $i \in \{1, 2, \ldots, m\}$. Next, we show that, for all $i \in \{1, 2, \ldots, m\}$,

$$H(Y_i|A, Y_1, \ldots, Y_{i-1}) \geq 1 - (16/\ln 2)\varepsilon^2.$$

(7)

To establish Eq. 7, it is useful to view the set $S$ as being generated during the execution of the $IP$ queries as follows. Initially $S$ is empty, and when the first $IP$ query is performed on some input $x$, $x$ is placed in $S$ with probability $\frac{1}{2} + \varepsilon$ and in $\overline{S}$ with probability $\frac{1}{2} - \varepsilon$. The inputs to subsequent $IP$ queries are also placed in either $S$ or $\overline{S}$ with an appropriate probability, which depends on how the inputs to previous queries are balanced between $S$ and $\overline{S}$. After the execution of the first $i - 1$ queries, the input to the $i$th query is placed in $S$ with probability

$$\frac{\left(\frac{1}{2} + \varepsilon\right)^{2n} - j}{2^n - (i - 1)};$$

(8)

where $j \in \{0, 1, \ldots, i - 1\}$ is the number of previous inputs to queries that have been placed in $S$. Using Eq. 3, the above probability can be shown to lie between $\frac{1}{2} - 2\varepsilon$ and $\frac{1}{2} + 2\varepsilon$. It follows that

$$H(Y_i|A, Y_1, \ldots, Y_{i-1}) \geq H\left(\frac{1}{2} + 2\varepsilon, \frac{1}{2} - 2\varepsilon\right)$$

$$= -(\frac{1}{2} + 2\varepsilon)\log(\frac{1}{2} + 2\varepsilon) - (\frac{1}{2} - 2\varepsilon)\log(\frac{1}{2} - 2\varepsilon)$$

$$\geq 1 - (16/\ln 2)\varepsilon^2,$$

(9)

establishing Eq. 7. Now, substituting the preceding inequalities into Eq. 6, we obtain

$$H(A|Y_1, \ldots, Y_m) \geq n - (16/\ln 2)m\varepsilon^2.$$

(10)

Intuitively, the $IP$ queries yield information about the value of $A$ in terms of their effect on the probability distribution of $A$ conditioned on the values of $Y_1, \ldots, Y_m$. Eq. 10 lower bounds the decrease in entropy possible.

From the conditions of the theorem, it can be assumed that, after the $IP$ queries, $2^{n/2}$ $EQ$ are performed. The algorithm succeeds with probability at least $\delta$ only if there exist $2^{n/2}$ elements of $\{0, 1\}^n$ whose total probability (conditioned on $Y_1, \ldots, Y_m$) is at least $\delta$. The maximum entropy that a distribution with this property can have is for a bi-level distribution, where $2^{n/2}$ elements of $\{0, 1\}^n$ each have probability $\delta/2^{n/2}$ and $2^n - 2^{n/2}$ elements each have probability $(1 - \delta)/(2^n - 2^{n/2})$. Therefore,

$$H(A|Y_1, \ldots, Y_m) \leq H\left(\frac{\delta}{2^{n/2}}, \ldots, \frac{\delta}{2^{n/2}}, \frac{1 - \delta}{2^{n/2}}, \ldots, \frac{1 - \delta}{2^{n/2}}, 0, \ldots, 0\right)$$

$$= H(\delta, 1 - \delta) + \delta \log(2^{n/2}) + (1 - \delta)\log(2^n - 2^{n/2})$$

$$< 1 + \delta n/2 + (1 - \delta)n$$

$$= n - \delta n/2 + 1.$$

(11)

Combining Eq. 10 with Eq. 11, yields $m > (\ln 2)(\delta n - 2)/(32\varepsilon^2) \in \Omega(\delta n/\varepsilon^2)$, as required.

We now provide definitions of $IP$ and $EQ$ queries in the quantum case in terms of unitary operations. We do this in a manner that is sufficiently general so that, whenever an implementation of a more general $IP$ or $EQ$ query is given as a general quantum circuit consisting of elementary quantum gates and measurements, a unitary query corresponding to our definition can be efficiently constructed from it.
Definition 3 A quantum inner product query (with bias $\varepsilon$) is a unitary transformation $U_{IP}$ on $n + m$ qubits, or its inverse $U_{IP}^\dagger$, such that $U_{IP}$ satisfies the following two properties:

1. If $x \in \{0, 1\}^n$ is chosen randomly according to the uniform distribution and the last qubit of $U_{IP}|x\rangle|0^m\rangle$ is measured, yielding the value $w \in \{0, 1\}$, then $\Pr[w = a \cdot x] \geq \frac{1}{2} + \varepsilon$.

2. For any $x \in \{0, 1\}^n$ and $y \in \{0, 1\}^m$, the state of the first $n$ qubits of $U_{IP}|x\rangle|y\rangle$ is $|x\rangle$.

The first property captures the fact that, taking a query to be a suitable application of $U_{IP}$ followed by a measurement of the last qubit, Eq. (1) is satisfied. Any implementation of a quantum circuit that produces an output that is $a \cdot x$ with probability on average $\frac{1}{2} + \varepsilon$ can be modified to consist of a unitary stage $U_{IP}$ followed by a measurement of one qubit. The second property is for technical convenience, and any unitary operation without this property can be converted to one that has this property, by first producing a copy of the classical basis state $|x\rangle$. Moreover, given a circuit implementing $U_{IP}$, it is easy to construct a circuit implementing $U_{IP}^\dagger$.

Definition 4 A quantum equivalence query is the unitary operation $U_{EQ}$ such that, for all $x \in \{0, 1\}^n$ and $b \in \{0, 1\}$,

$$U_{EQ}|x\rangle|b\rangle = \begin{cases} |x\rangle|\overline{b}\rangle & \text{if } x = a \\ |x\rangle|b\rangle & \text{if } x \neq a, \end{cases} \quad (12)$$

where $\overline{b} = -b$.

For the quantum GL problem, $a \in \{0, 1\}^n$ and information about $a$ in available only from quantum $IP$ and $EQ$ queries and the goal is to determine $a$. We can now state and prove the result about quantum algorithms for the GL problem (which is similar to a result in [3] in a different context).

Theorem 2 There exists a quantum algorithm solving the GL problem with constant probability using $O(1/\varepsilon)$ $U_{IP}$, $U_{IP}^\dagger$ and $U_{EQ}$ queries in total. Also, the number of auxiliary qubit operations used by the procedure is $O(n/\varepsilon)$.

Proof: The proof is by a combination of two techniques: the algorithm in [3] for the exact case (i.e., when $\varepsilon = \frac{1}{2}$), which is shown to be adaptable to “noisy” data in [8] (with a slightly different noise model than the one that arises here); and amplitude amplification [6].

Since $U_{IP}$ applied to $|x\rangle|y\rangle$ has no net effect on its first $n$ input qubits, for each $x \in \{0, 1\}^n$,

$$U_{IP}|x\rangle|0^m\rangle = |x\rangle (\alpha_x|v_x\rangle|a \cdot x\rangle + \beta_x|w_x\rangle|\overline{a \cdot x}\rangle), \quad (13)$$

where $\alpha_x$ and $\beta_x$ are nonnegative real numbers, and $|v_x\rangle$ and $|w_x\rangle$ are $m - 1$ qubit quantum states. If the last qubit of $U_{IP}|x\rangle|0^m\rangle$ is measured then the result is: $a \cdot x$ with probability $\alpha_x^2 + \beta_x^2$, and $\overline{a \cdot x}$ with probability $\overline{\beta_x^2}$. Therefore, since, for a random uniformly distributed $x \in \{0, 1\}^n$, measuring the last qubit of $U_{IP}|x\rangle|0^m\rangle$ yields $a \cdot x$ with probability at least $\frac{1}{2} + \varepsilon$, it follows that

$$\frac{1}{2n} \sum_{x \in \{0, 1\}^n} \alpha_x^2 \geq \frac{1}{2} + \varepsilon \quad (14)$$

$$\frac{1}{2n} \sum_{x \in \{0, 1\}^n} \beta_x^2 \leq \frac{1}{2} - \varepsilon. \quad (15)$$
Now, consider the quantum circuit $C$ in Figure 1.

Figure 1: Quantum circuit $C$.

We will begin by showing that $\langle a, 0^m, 1 | C | 0^n, 0^m, 0 \rangle$ is real-valued and

$$\langle a, 0^m, 1 | C | 0^n, 0^m, 0 \rangle \geq 2\varepsilon,$$  \hspace{1cm} (16)

which intuitively can be viewed as an indication of the progress that $C$ makes towards finding the string $a$. To establish Eq. (16), note that the operation $C$ can be decomposed into the following five operations:

1. Operation $C_1$: Apply $H$ to each of the first $n$ qubits, and a NOT operation to the last qubit.
2. Operation $C_2$: Apply $U_{IP}$ to the first $n + m$ qubits.
3. Operation $C_3$: Apply a controlled-$Z$ to the last two qubits.
4. Operation $C_4$: Apply $U_{IP}^\dagger$ to the first $n + m$ qubits.
5. Operation $C_5$: Apply $H$ to each of the first $n$ qubits.

Since $\langle a, 0^m, 1 | C | 0^n, 0^m, 0 \rangle = \langle a, 0^m, 1 | C_5 C_4 C_3 C_2 C_1 | 0^n, 0^m, 0 \rangle$, the quantity $\langle a, 0^m, 1 | C | 0^n, 0^m, 0 \rangle$ is the inner product between state $C_3 C_2 C_1 | 0^n \rangle | 0^m \rangle | 0 \rangle$ and state $C_4^\dagger C_5^\dagger | a \rangle | 0^m \rangle | 1 \rangle$. These states are

$$C_3 C_2 C_1 | 0^n \rangle | 0^m \rangle | 0 \rangle = C_3 C_2 \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle |0^m\rangle |1\rangle$$

$$= C_3 \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle (\alpha_x |v_x\rangle |a \cdot x\rangle + \beta_x |w_x\rangle |a \cdot x\rangle) \langle 1 |$$

$$= \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle (\alpha_x (-1)^a \cdot x |v_x\rangle |a \cdot x\rangle + \beta_x (-1)^a \cdot x |w_x\rangle |a \cdot x\rangle) \langle 1 |$$

$$= \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^a \cdot x |x\rangle (\alpha_x |v_x\rangle |a \cdot x\rangle - \beta_x |w_x\rangle |a \cdot x\rangle) \langle 1 |$$  \hspace{1cm} (17)
and

\[
C_4^\dagger C_5^\dagger |a\rangle |0^m\rangle |1\rangle = C_4^\dagger \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^a \cdot x |x\rangle |0^m\rangle |1\rangle \\
= \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^a \cdot x (\alpha_x |v_x\rangle |a \cdot x\rangle + \beta_x |w_x\rangle |a \cdot \overline{x}\rangle) |1\rangle,
\]

(18)

It follows from Eq. 17 and Eq. 18 (and using the fact that \langle x|y \rangle = 0 whenever \( x \neq y \)) that

\[
\langle a, 0^m, 1 | C | 0^n, 0^m, 0 \rangle = \frac{1}{2^n} \sum_{x \in \{0,1\}^n} (\alpha_x^2 - \beta_x^2) \\
\geq \left( \frac{1}{2} + \varepsilon \right) - \left( \frac{1}{2} - \varepsilon \right) \\
= 2\varepsilon,
\]

(19)

which establishes Eq. 14.

Note that Eq. 14 implies that, if \( C \) is executed on input \( |0^n\rangle |0^m\rangle |0 \rangle \) and the result is measured in the classical basis, then the first \( n \) bits of the result will be \( a \) with probability at least \(|\langle a, 0^m, 1 | C | 0^n, 0^m, 0 \rangle|^2 \geq 4\varepsilon^2 \). Therefore, if this process is repeated \( O(1/\varepsilon^2) \) times, checking each result with an \( EQ \) query, then \( a \) will be found with constant probability. A more efficient way of finding the value of \( a \) is to use amplitude amplification \[5, 15, 6\] using the transformation \( C \) and its inverse \( C^\dagger \) in combination with \( EQ \) queries. The procedure is to compute (for various values of \( k \))

\[
(-CU_0C^\dagger U_{EQ})^k C | 0^n, 0^m, 0 \rangle
\]

(20)

(where \( U_0 = I - 2|0^n, 0^m, 0 \rangle \langle 0^n, 0^m, 0 | \)), measure the state, and perform an \( EQ \) query on the result. Such a computation consists of \( O(k) \) \( U_{IP}, U_{IP}^\dagger \), and \( U_{EQ} \) queries. As shown in [6], if this is carried out for a suitably generated sequence of values of \( k \), the expected total number of executions of \( C, C^\dagger, \) and \( U_{EQ} \) until a successful \( EQ \) query occurs is \( O(1/\varepsilon) \). This implies that \( O(1/\varepsilon) \) \( U_{IP}, U_{IP}^\dagger \), and \( U_{EQ} \) are sufficient to succeed with constant probability.

\[ \]

3 Hard-predicates from one-way permutations

In this section, we give definitions pertaining to one-way permutations and hard-predicates (classical and quantum versions) and investigate the complexity of the reduction of Goldreich and Levin \[13\] from the former to the latter.\footnote{The reduction makes sense for functions that are not permutations, but we restrict attention to permutations for simplicity.}

In the definitions below, when we refer to the \textit{size} of a classical [quantum] circuit, it is understood to be relative to a suitable set of gates on one and two bits [qubits]. Quantum circuits compute unitary transformations on quantum states; however, they can also be adapted to take classical data as input and produce classical data as output. For a quantum circuit \( C \) acting on \( m \) qubits, and \( x \in \{0,1\}^n \) (for \( n \leq m \)), let \( C_k(x) \) \((k \in \{1, \ldots , m\})\), denote the result of measuring the first \( k \) qubits of \( C|x\rangle |0^{m-n}\rangle \) in the classical basis. The subscript \( k \) may be omitted when the value of \( k \) is clear from the context.
Intuitively, a quantum one-way permutation $f$ on $n$ bits is easy to compute in the forward direction but is hard to invert\footnote{The reversibility of quantum computations does not exclude this possibility \cite{4}.}. For the former property, the standard requirement is that $f$ be computable by a uniform circuit of size $n^{O(1)}$ (though it is also possible to impose other upper bounds on the uniform circuit size). To quantify the latter property, it is helpful to first make the following definition.

**Definition 5** A permutation $f : \{0,1\}^n \to \{0,1\}^n$ is classically [quantumly] $(\delta, T)$-hard to invert if there is no classical [quantum] circuit $C$ of size $T$ such that $\Pr_a[C(f(a)) = a] \geq \delta$.

Now the standard requirement for the hard-to-invert condition is that $f$ is $(\delta, T)$-hard to invert for all $\delta \in 1/n^{O(1)}$ and $T \in n^{O(1)}$ (again, other bounds can be imposed). It should be noted that, although it may be hard to determine $a$ from $f(a)$, it may not be hard to extract partial information about $a$ from $f(a)$. For example, it is conceivable for a one-way permutation $f$ to have the property that half of the bits of $a$ can be efficiently determined exactly from $f(a)$. It is also conceivable that each individual bit of $a$ is efficiently predictable from $f(a)$ with probability $\frac{1}{2}$. The idea behind a hard-predicate \cite{4} is to concentrate the information that a one-way function “hides” about its input into a single bit. Intuitively, $h : \{0,1\}^n \to \{0,1\}$ is a hard-predicate of $f$ if, given $a \in \{0,1\}^n$, it is easy to compute $h(a)$; whereas, given $f(a)$ for randomly chosen $a \in \{0,1\}^n$, it is hard to predict the value of the bit $h(a)$ with probability significantly better than $\frac{1}{2}$. One natural way of quantifying how well a circuit predicts the value of $h$ from the value of $f$ is by the amount that $\Pr_a[C(f(a)) = h(a)]$ exceeds $\frac{1}{2}$.

The hard-predicate defined in \cite{13} is

$$h(y, x) = y \cdot x,$$

(\text{the inner product modulo two of $x$ and $y$}), for $(y, x) \in \{0,1\}^n \times \{0,1\}^n$. This is not a hard-predicate of $f$, but for a slightly modified version of $f$, as given in the following definition.

**Definition 6** For a permutation $f : \{0,1\}^n \to \{0,1\}^n$, let $\tilde{f}$ denote the permutation $\tilde{f} : \{0,1\}^n \times \{0,1\}^n \to \{0,1\}^n \times \{0,1\}^n$ defined as

$$\tilde{f}(y, x) = (f(y), x),$$

for all $(y, x) \in \{0,1\}^n \times \{0,1\}^n$.

Note that the cost of computing [inverting] $\tilde{f}$ is essentially the same as the cost of computing [inverting] $f$. Goldreich and Levin showed that if $f$ is one-way then $h$ is hard to predict from $\tilde{f}$. Instead of quantifying how well a circuit predicts $h$ from $\tilde{f}$ as the amount by which $\Pr_{y, x}[C(\tilde{f}(y, x)) = h(y, x)]$ exceeds $\frac{1}{2}$, we adopt a slightly more complicated definition. This definition is related to the above, but is better suited for expressing the results in this section.

**Definition 7** A circuit $C$ $(\delta, \varepsilon)$-predicts $h$ from $\tilde{f}$ if

$$\Pr_y[\Pr_x[C(\tilde{f}(y, x)) = h(y, x)] \geq \frac{1}{2} + \varepsilon] \geq \delta.$$

To explain Eq. \ref{eq:4} in words, call $y \in \{0,1\}^n$ \emph{\varepsilon-good} if $\Pr_x[C(\tilde{f}(y, x)) = h(y, x)] \geq \frac{1}{2} + \varepsilon$ for that value of $y$. Then then Eq. \ref{eq:4} is equivalent to saying that $\Pr_y[y \text{ is \varepsilon-good}] \geq \delta$.\footnote{The reversibility of quantum computations does not exclude this possibility \cite{4}.}
The following lemma, which relates the two measures of prediction, is straightforward to prove by an averaging argument.

**Lemma 3** If \( \Pr_{y,x}[G(\tilde{f}(y,x)) = h(y,x)] \geq \frac{1}{2} + \varepsilon \) then \( G(\varepsilon/(1-\varepsilon), \varepsilon/2) \)-predicts \( h \) from \( \tilde{f} \).

Note that, if \( \Pr_{y,x}[G(\tilde{f}(y,x)) = h(y,x)] \geq \frac{1}{2} + 1/n^{O(1)} \) then \( G(1/n^{O(1)}, 1/n^{O(1)}) \)-predicts \( h \) from \( \tilde{f} \).

The classical Goldreich-Levin Theorem can be stated as follows.

**Theorem 4** ([13, 14]) If \( f : \{0,1\}^n \rightarrow \{0,1\}^n \) is classically \((\delta/2, T)\)-hard to invert then any classical circuit that \((\delta, \varepsilon)\)-predicts \( h \) from \( \tilde{f} \) must have size \( \Omega(T\varepsilon^2/n) \).

The proof of this theorem is essentially a reduction from the problem of inverting \( f \) to the problem of \((\delta, \varepsilon)\)-predicting \( h \). One begins by assuming that a circuit \( G \) of size \( o(T\varepsilon^2/n) \) \((\delta, \varepsilon)\)-predicts \( h \) from \( \tilde{f} \) and then shows that, by making \( O(n/\varepsilon^2) \) calls to both \( G \) and \( f \) (plus some additional computations), \( f \) can be inverted with probability \( \delta/2 \) [14]. The total running time of the inversion procedure is \( o((n/\varepsilon^2)(T\varepsilon^2/n)) = o(T) \), contradicting the fact that \( f \) is \((\delta/2, T)\)-hard to invert.

Our quantum version of the Goldreich-Levin Theorem is the following.

**Theorem 5** If \( f : \{0,1\}^n \rightarrow \{0,1\}^n \) is quantumly \((\delta/2, T)\)-hard to invert then any quantum circuit that \((\delta, \varepsilon)\)-predicts \( h \) from \( \tilde{f} \) must have size \( \Omega(T\varepsilon) \).

**Proof:** As in the classical case, the proof is essentially a reduction from the problem of inverting \( f \) to the problem of \((\delta, \varepsilon)\)-predicting \( h \). Let \( b = f(a) \) be an input instance—the goal is to determine \( a \) from \( b \). We will show how to simulate \( EQ \) and \( IP \) queries in this setting and then apply the bounds in Theorem 2. It is easy to simulate an \( EQ \) query (relative to \( a \)) by making one call to \( f \) and checking if the result is \( b \). Suppose that there exists a circuit \( G \) of size \( o(T\varepsilon) \) that \((\delta, \varepsilon)\)-predicts \( h \) from \( \tilde{f} \). Thus, \( \Pr_y[\Pr_x[G(\tilde{f}(y,x)) = h(y,x)] \geq \frac{1}{2} + \varepsilon] \geq \delta \). Note that, with probability at least \( \delta \), \( a \) is \( \varepsilon \)-good, in the sense that \( \Pr_x[G(\tilde{f}(a,x)) = h(a,x)] \geq \frac{1}{2} + \varepsilon \). When \( a \) is \( \varepsilon \)-good, computing \( G(\tilde{f}(a,x)) = G(b,x) \) is simulating an \( IP \) query for \( x \) (relative to \( a \)). It follows from Theorem 2 that \( a \) can be computed with circuit-size \( o(1/\varepsilon)(T\varepsilon) = o(T) \) with success probability at least \( \delta/2 \) (where \( 1/2 \) is the success probability of the algorithm that finds \( a \) when \( a \) is \( \varepsilon \)-good and \( \delta \) is the probability that \( a \) is \( \varepsilon \)-good to begin with). This contradicts the \((\delta/2, T)\)-hardness of inverting \( f \), thus \( G \) cannot \((\delta, \varepsilon)\)-predict \( h \) from \( \tilde{f} \) and be of size \( o(T\varepsilon) \).

To conclude this section, we give a proof that Theorem 4 cannot be improved quantitatively assuming that it follows the structure of making calls to \( f \) and to an algorithm \( G \) that \((\delta, \varepsilon)\)-predicts \( h \) from \( \tilde{f} \). More precisely, the setting is as follows. For a permutation \( f : \{0,1\}^n \rightarrow \{0,1\}^n \), information is available from two types of black-box queries: \( f \)-queries that evaluates \( f \); and \( G \)-queries that \((\delta, \varepsilon)\)-predict \( h \) from \( \tilde{f} \). More precisely, a \( G \)-query has the property that \( \Pr_y[\Pr_x[G(\tilde{f}(y,x)) = h(y,x)] \geq \frac{1}{2} + \varepsilon] \geq \delta \). A problem instance is \( b \in \{0,1\}^n \) (where \( b = f(a) \) for a random \( a \in \{0,1\}^n \)) and the availability of \( f \)-queries and \( G \)-queries. The goal is to determine \( a \) with probability \( \delta/2 \) (say). Let us refer to this as the \( GL^* \) problem (related to but different from the GL problem defined in Section 2). From the proof of Theorem 3, the classical \( GL^* \) problem can be solved with \( O(n/\varepsilon^2) \) \( f \)-queries and \( G \)-queries (and \( O(n^2/\varepsilon^2) \) auxiliary operations [14]). From the proof of Theorem 3, a quantum version of the \( GL^* \) problem can be solved with only \( O(1/\varepsilon) \) \( f \)-queries and \( G \)-queries (and \( O(n/\varepsilon) \) auxiliary operations). The next theorem essentially implies that the dilution factor \( n/\varepsilon^2 \) in Theorem 4 cannot be reduced for a reasonable range of values of \( \varepsilon \).
Theorem 6 The classical GL\(^*\) problem requires either \(\Omega(2^{n/2})\) \(f\)-queries or \(\Omega(n/\varepsilon^2)\) \(G\)-queries, whenever \(\varepsilon \geq \sqrt{n/2^{-n/4}}\).

Proof: The idea behind the proof is show that, starting with an algorithm that solves the GL\(^*\) problem using \(T_f\) \(f\)-queries and \(T_G\) \(G\)-queries, it is possible to simulate each \(f\)-query with one \(EQ\) query and to simulate each \(G\)-query with one \(IP\) query and one \(EQ\) query. The result is an algorithm that solves the GL problem defined in Section 2 with \(T_G\) \(IP\) queries and \(T_f + T_G\) \(EQ\) queries. Then, applying the bound in Theorem 4, yields the required lower bounds.

By a basic game-theoretic argument \([20]\), it suffices to consider deterministic algorithms, where the input data—embodied by

\[\text{the queries. Then, applying the bound in Theorem 1, yields the required lower bounds.}\]

The idea behind the proof is show that, starting with an algorithm that solves the GL\(^*\) problem using \(T_f\) \(f\)-queries and \(T_G\) \(G\)-queries, it is possible to simulate each \(f\)-query with one \(EQ\) query and to simulate each \(G\)-query with one \(IP\) query and one \(EQ\) query. The result is an algorithm that solves the GL problem defined in Section 2 with \(T_G\) \(IP\) queries and \(T_f + T_G\) \(EQ\) queries. Then, applying the bound in Theorem 4, yields the required lower bounds.

The idea behind the proof is show that, starting with an algorithm that solves the GL\(^*\) problem using \(T_f\) \(f\)-queries and \(T_G\) \(G\)-queries, it is possible to simulate each \(f\)-query with one \(EQ\) query and to simulate each \(G\)-query with one \(IP\) query and one \(EQ\) query. The result is an algorithm that solves the GL problem defined in Section 2 with \(T_G\) \(IP\) queries and \(T_f + T_G\) \(EQ\) queries. Then, applying the bound in Theorem 4, yields the required lower bounds.

The idea behind the proof is show that, starting with an algorithm that solves the GL\(^*\) problem using \(T_f\) \(f\)-queries and \(T_G\) \(G\)-queries, it is possible to simulate each \(f\)-query with one \(EQ\) query and to simulate each \(G\)-query with one \(IP\) query and one \(EQ\) query. The result is an algorithm that solves the GL problem defined in Section 2 with \(T_G\) \(IP\) queries and \(T_f + T_G\) \(EQ\) queries. Then, applying the bound in Theorem 4, yields the required lower bounds.

The idea behind the proof is show that, starting with an algorithm that solves the GL\(^*\) problem using \(T_f\) \(f\)-queries and \(T_G\) \(G\)-queries, it is possible to simulate each \(f\)-query with one \(EQ\) query and to simulate each \(G\)-query with one \(IP\) query and one \(EQ\) query. The result is an algorithm that solves the GL problem defined in Section 2 with \(T_G\) \(IP\) queries and \(T_f + T_G\) \(EQ\) queries. Then, applying the bound in Theorem 4, yields the required lower bounds.
4 Quantum bit commitment from quantum one-way permutations

In this section, we show how to use the quantum Goldreich-Levin Theorem to construct a quantum bit commitment scheme from a quantum one-way permutation.

Definition 8 A permutation \( f : \{0,1\}^n \rightarrow \{0,1\}^n \) is a quantum one-way permutation if:

- There is a uniform quantum circuit of size \( n^{O(1)} \) that computes \( f(x) \) from \( x \).
- \( f \) is quantumly \((\delta,T)\)-hard to invert for any \( \delta \in 1/n^{O(1)} \) and \( T \in n^{O(1)} \).

Theorem 7 If there exists a quantum one-way permutation \( f : \{0,1\}^n \rightarrow \{0,1\}^n \) then there exists a bit [or qubit] commitment scheme that is perfectly binding and computationally concealing, in the sense that the committed bit cannot be predicted with probability \( \frac{1}{2} + 1/n^{O(1)} \) by a circuit of size \( n^{O(1)} \).

Proof: From Theorem 5, it is straightforward to construct a quantum bit commitment scheme from a quantum one-way permutation.

\[ \text{Bit-commit} \]

Let \( \psi \) be the qubit to commit to. Alice chooses \( a, x \in \{0,1\}^n \) randomly, and sets \( c = z \oplus h(a, x) \). Alice computes \( b = f(a) \) and sends \((b, x, c) \) to Bob.

\[ \text{Bit-decommit} \]

Alice sends \( a \) to Bob. Bob checks if \( f(a) = b \) and rejects if this is not the case. Otherwise, Bob accepts and computes \( c \oplus h(a, x) \) as the bit.

Since \( f \) is a permutation there is at most one classical value of \( a \) that is an acceptable decommitment of Alice’s bit. This implies that the scheme is perfectly binding to Alice. Note that the model could be relaxed to permit Alice to send quantum data to Bob, by adjusting Bob’s protocol to immediately perform a measurement (in the classical basis) on any data that he receives from Alice. There would be no advantage to Alice—she could not somehow “commit to more than one value” by sending commitments in superposition. This is because the adjusted protocol is equivalent to one where Alice performs the measurement herself on any data before sending it to Bob.

Theorem 5 implies that the scheme is also computationally concealing, since any \( n^{O(1)} \)-size circuit that enables Bob to guess \( z \) from \((b, x, c)\) with probability \( \frac{1}{2} + 1/n^{O(1)} \) can be converted to a \( n^{O(1)} \)-size circuit that inverts \( f \) with probability \( 1/n^{O(1)} \), violating the fact that \( f \) is one-way.

Finally, we explain how a qubit commitment scheme can be constructed using some of the ideas in [2]. Recall the standard notation for the Pauli matrices:

\[
X = \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \text{and} \quad Z = \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.
\] (24)

\[ \text{Qubit-commit} \]

Let \( |\psi\rangle \) be the qubit to commit to. Alice chooses \( a_1, a_2, x_1, x_2 \in \{0,1\}^n \) randomly, and constructs the state \( |\psi'\rangle = X^{h(a_1,x_1)}Z^{h(a_2,x_2)}|\psi\rangle \) and also computes \( b_1 = f(a_1) \) and \( b_2 = f(a_2) \). Alice sends \((|\psi'\rangle, b_1, b_2, x_1, x_2) \) to Bob.

\[ \text{Qubit-decommit} \]

Alice sends \( a_1, a_2 \) to Bob. Bob checks if \( f(a_1) = b_1 \) and \( f(a_2) = b_2 \), rejecting if this is not the case. Otherwise, Bob accepts and computes \( Z^{h(a_2,x_2)}X^{h(a_1,x_1)}|\psi'\rangle \) as the qubit.
Clearly, the scheme is perfectly binding. Intuitively, the scheme is computationally concealing, because \( h(a_1, x_1) \) and \( h(a_2, x_2) \) “look random” to Bob. If Bob can use his information to efficiently significantly distinguish between the qubit that he receives from Alice in the commitment stage and a totally mixed state (density matrix \( \frac{1}{2}I \)) then this procedure can be adapted to distinguish between the pair of bits \( r_1 = h(a_1, x_1) \) and \( r_2 = h(a_2, x_2) \) and a pair of truly random bits, which would lead to a procedure that violated the result proven in Theorem 5.

Acknowledgments

We would like to thank Paul Dumais, Peter Høyer, Dominic Mayers, and Louis Salvail for helpful discussions.

References

[1] A. Ambainis, M. Mosca, A. Tapp, R. de Wolf, “Private quantum channels”, Proc. 41st Ann. IEEE Symp. on Foundations of Computer Science (FOCS ’00), pp. 547–553, 2000.

[2] M. Bellare, “The Goldreich-Levin Theorem”, Manuscript, 1999. (Available at http://www-cse.ucsd.edu/users/mihir/)

[3] E. Bernstein and U. V. Vazirani, “Quantum complexity theory”, SIAM J. on Comput., Vol. 26, No. 5, pp. 1411–1473, 1997.

[4] M. Blum and S. Micali, “How to generate cryptographically strong sequences of pseudo-random bits”, SIAM J. on Comput., Vol. 13, No. 4, pp. 850–864, 1984.

[5] G. Brassard and P. Høyer, “An exact quantum polynomial-time algorithm for Simon’s problem”, Proc. Fifth Israeli Symp. on Theory of Computing and Systems, pp. 12–23, 1997.

[6] G. Brassard, P. Høyer, M. Mosca, A. Tapp, “Quantum amplitude amplification and estimation”, To appear in Quantum Computation and Quantum Information: A Millennium Volume, AMS Contemporary Mathematics Volume. Available on the LANL preprint archive as quant-ph/0005055, 2000.

[7] H. F. Chau and H.-K. Lo, “One way functions in reversible computations”, Cryptologia, Vol. 21, No. 2, pp. 139–148, 1997.

[8] R. Cleve, W. van Dam, M. Nielsen, and A. Tapp, “Quantum entanglement and the communication complexity of the inner product function”, Proc. of the First NASA International Conf. on Quantum Computing and Quantum Communications, Colin P. Williams (Ed.), Lecture Notes in Computer Science 1509, Springer-Verlag, pp. 61-74, 1999.

[9] T. M. Cover and J. A. Thomas, Elements of Information Theory, John Wiley and Sons, 1991.

[10] C. Crépeau, F. Légaré and L. Salvail, “How to convert the flavor of a quantum bit commitment”, to appear in Advances in Cryptology — EUROCRYPT 2001.
[11] P. Dumais, D. Mayers, and L. Salvail, “Perfectly concealing quantum bit commitment from any one-way permutation”, *Advances in Cryptology — EUROCRYPT 2000*, B. Preneel (Ed.), Lecture Notes in Computer Science 1807, Springer-Verlag, pp. 300–315, 2000.

[12] M. R. Garey and D. S. Johnson, *Computers and Intractability: A Guide to the Theory of NP-Completeness*, W. H. Freeman & Co., 1979.

[13] O. Goldreich and L. Levin, “Hard-core predicates for any one-way function”, *Proc. 21th Ann. ACM Symp. on Theory of Computing (STOC ’89)*, pp. 25–32, 1989.

[14] O. Goldreich, *Modern Cryptography, Probabilistic Proofs and Pseudo-Randomness*, Springer, 1999.

[15] L. K. Grover, “A fast quantum mechanical algorithm for database search”, *Proc. 28th Ann. ACM Symp. on Theory of Computing (STOC ’96)*, pp. 212–219, 1996.

[16] H.-K. Lo and H. F. Chau, “Is quantum bit commitment really possible?”, *Phys. Rev. Lett.*, Vol. 78, No. 17, pp. 3410–3413, 1997.

[17] D. Mayers, “Unconditionally secure bit commitment is impossible”, *Phys. Rev. Lett.*, Vol. 78, No. 17, pp. 3414–3417, 1997.

[18] P. W. Shor, “Polynomial-time algorithms for prime factorization and discrete logarithms on a quantum computer”, *SIAM J. on Computing*, Vol. 26, No. 5, pp. 1484–1509, 1997.

[19] B. M. Terhal and J. A. Smolin, “Single quantum querying of a database”, *Phys. Rev. A*, Vol. 58, No. 3, pp. 1822–1826, 1998.

[20] A. C.-C. Yao, “Lower bounds by probabilistic arguments”, *Proc. 24th Ann. IEEE Symp. on Foundations of Computer Science (FOCS ’83)*, pp. 420–428, 1983.