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Title: Nuclear spin hyperpolarization with ansa-aminoboranes: a metal-free perspective for parahydrogen-induced polarization

This paper reports a series of unimolecular metal-free frustrated Lewis pairs based on an ansa-aminoborane (AAB) moiety, allowing observation of nuclear spin hyperpolarization effects upon interaction with parahydrogen. Hyperpolarization was observed for all AABs differing in boryl site substituent (–H, –Ph, –o-iPr-Ph, –Mes) in a wide temperature range. A theoretical analysis revealed the roles of chemical exchange, chemical equilibrium and spin dynamics in terms of the key parameters, providing a foundation for design of metal-free tweezer-like molecules for parahydrogen-induced polarization.

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Nuclear spin hyperpolarization with \textit{ansa}-aminoboranes: a metal-free perspective for parahydrogen-induced polarization\dagger

Vladimir V. Zhivonitko,\textsuperscript{a}b Kristina Sorochkina,\textsuperscript{c} Konstantin Chernichenko,\textsuperscript{c} Bianka Kótai,\textsuperscript{d} Tamás Földes,\textsuperscript{d} Imre Pápai,\textsuperscript{d} Ville-Veikko Telkki,\textsuperscript{e} Timo Repo\textsuperscript{c} and Igor Koptyug\textsuperscript{ab}

The parahydrogen-induced polarization (PHIP) phenomenon, observed when parahydrogen is used in H\textsubscript{2} addition processes, provides a means for substantial NMR signal enhancements and mechanistic studies of chemical reactions. Commonly, noble metal complexes are used for parahydrogen activation, whereas metal-free activation is rare. Herein, we report a series of unimolecular metal-free frustrated Lewis pairs based on an \textit{ansa}-aminoborane (AAB) moiety in the context of PHIP. These molecules, which have a “molecular tweezers” structure, differ in their substituents at the boryl site (\(-\text{H}, \sim\text{Ph}, \sim\text{o-iPr-Ph}, \text{and} \sim\text{Mes}\)). PHIP effects were observed for all the AABs after exposing their solutions to parahydrogen in a wide temperature range, and experimental measurements of their kinetic and thermodynamic parameters were performed. A theoretical analysis of their nuclear spin polarization effects is presented, and the roles of chemical exchange, chemical equilibrium and spin dynamics are discussed in terms of the key dimensionless parameters. The analysis allowed us to formulate the prerequisites for achieving strong polarization effects with AAB molecules, which can be applied for further design of efficient metal-free tweezers-like molecules for PHIP. Mechanistic (chemical and physical) aspects of the observed effects are discussed in detail. In addition, we performed quantum chemical calculations, which confirmed that the J-coupling between the parahydrogen-originated protons in AAB–H\textsubscript{2} molecules is mediated through dihydrogen bonding.

1. Introduction

The high reactivity of frustrated Lewis pairs (FLPs) was discovered only recently,\textsuperscript{1,2} inspiring numerous brilliant studies in this field over the last few years.\textsuperscript{3–12} Many small molecules have been shown to be activated by FLPs\textsuperscript{6,12} including dihydrogen (H\textsubscript{2}). In addition to the apparent applications in catalysis,\textsuperscript{13–24} addition of H\textsubscript{2} to FLPs may be useful for hyperpolarizing nuclear spins when parahydrogen, one of the nuclear spin isomers of H\textsubscript{2}, is used as a reagent. Due to the symmetry of the H\textsubscript{2} molecule, the correlation of its nuclear spins survives for a long time and can be transformed into observable nuclear spin hyperpolarization upon a chemical transformation of H\textsubscript{2} molecules if this transformation causes the two hydrogens to be magnetically inequivalent. At the same time, the H\textsubscript{2} activation mechanism must be pairwise, meaning that both hydrogens should follow each other throughout all elementary reaction steps and end up in the same product molecule. Because of this prerequisite, only certain catalytic systems are applicable in this hyperpolarization technique, which is referred to in the literature as PASADENA,\textsuperscript{25} ALTADENA,\textsuperscript{26} or, generally, parahydrogen-induced polarization (PHIP).\textsuperscript{27} To date, PHIP mediated by FLPs was demonstrated only in a few examples as a proof-of-principle.\textsuperscript{28,29} In particular, it has been shown that the reversibility of the H\textsubscript{2} interaction with FLPs allows one to achieve reproducible nuclear spin hyperpolarization each time the parahydrogen flux is introduced to a QCAT ansa-aminoborane.
solution (Scheme 1), imitating a kind of nuclear spin polarization “pumping”.28

In addition to the production of nuclear spin hyperpolarization for substantial signal enhancement in NMR and MRI, PHIP can shed light on the mechanistic features of molecular activation. Particularly, parahydrogen provides a unique type of labeling which is sensitive to the pairwise mechanism. At the same time, progress in understanding the involved mechanisms can provide a means for the development of new efficient FLP systems for PHIP. In particular, determination and optimization of the kinetic parameters of the activation are also crucial for maximizing PHIP. Unlike the spin correlation in parahydrogen, the lifetime of nuclear spin polarization in FLP–H₂ adducts is relatively short, which strongly diminishes the achievable polarization levels if the kinetics of the hydrogen addition is slow. This actually means that high addition rates compared to the spin relaxation rates are favorable for the observation of high polarization levels.

Recently, we reported that ansa-aminoboranes (AABs), N–B based FLPs, demonstrate PHIP effects, indicating that these compounds possess the required properties.28 The chelating arrangement of the N and B reactive centers in AABs favors rapid H₂ addition dynamics. Importantly, the intrinsic feature of AABs is the presence of a dihydrogen NH· · · HB bond in the AAB–H₂ adducts. This bond is evident from X-ray and neutron diffraction studies29 and from computational results. Thus, AABs are referred to in the literature as “molecular tweezers” that stretch, but do not completely split, H₂ molecules. As shown below, the dihydrogen bond enables the observation of PHIP. Apparently, with further development, AABs may have great prospects in PHIP production. Namely, an appropriate molecular design could allow one to use close-to-ambient temperatures and, as an ultimate goal, to design moisture-tolerant AABs for in vivo MRI applications. Fundamentally, molecular activation with these compounds is of exceptional interest for metal-free catalysis as a new area with major potential for significant future advances.

Switching from the originally developed AABs, which were based on the 2-boryl-benzylamine core, to 2-aminophenylboranes provided more favorable thermodynamics of H₂ addition.31 Due to its unique properties, which enable unprecedented catalytic hydrogenations31 and C–H borylations,32,33 the 2-aminophenylboryl core allowed us to design novel AABs with substituents other than C₆F₅ at the boryl site. Particularly, we have recently reported that BorylCAT, the AAB with the smallest boryl site (–BH₂), also demonstrates a reversible dihydrogen addition at ambient temperatures.34 In the present work, we prepared several analogues of BorylCAT which contain various aryl substituents (phenyl, mesityl, 2-(isopropyl)phenyl) at the boryl site (Scheme 1). These di(aryl)boryl-substituted AABs, together with BorylCAT, constitute a series of isostuctural compounds with varied electronic and, especially, steric properties of their Lewis acidic sites. We found that all these AABs add hydrogen with nearly neutral thermodynamics at close to ambient conditions. Moreover, we show that PHIP effects can be observed for these compounds under similar conditions. The influences of reaction kinetics, relaxation and scalar spin coupling parameters on the intensity of the PHIP signal is discussed in detail. Based on computational analysis of the BorylCAT–H₂ adduct, it is shown that the scalar spin coupling between the two protons originating from H₂ is the result of dihydrogen bonding and that borate (–BH₄) fragment rotation (libration) occurs.

2. Experimental

2.1 PHIP experiments

PHIP NMR experiments were performed on a Bruker AV 300 SB NMR spectrometer operating at 300 MHz ¹H resonance frequency and equipped with a broad-band 10 mm RF probe. The standard temperature control unit of the NMR spectrometer was used to maintain the desired temperature of the sample. The sample temperature was controlled by programatically adjusted heating of a cold N₂ flow from a liquid nitrogen tank.

Commercially available high-purity H₂ gas was used for producing parahydrogen-enriched H₂, referred to below as “para-H₂”. In this enrichment, H₂ from a gas cylinder was supplied through a spiral-shaped copper tube immersed in liquid N₂ at 77 K. The ortho–para conversion catalyst (FeO(OH)) was packed at the
end of the spiral. As cooled H₂ flowed through the cold catalyst, para-hydrogen-enriched H₂ with a 1:1 ortho:para ratio was produced for further use in the experiments.

In a typical workflow, a 5 mm sample tube (Wilmad) containing a ~0.05 M solution of AAB in degassed dry methylene chloride-d₂ was placed inside the NMR magnet, and para-H₂ was obtained as described above, was bubbled through the solution for 10 s. Then, the parahydrogen flow was abruptly terminated and the NMR experiment was started. ¹⁴N decoupling (WALTZ16) was used for the acquisitions.

The 5 mm tube was fixed in a heavy-wall 10 mm tube including a CCl₄ jacket, stabilizing the temperature in the solution (Fig. 1). para-H₂ was supplied to the bottom of the sample tube through 1/32” PTFE tubing under a pressure of 3.2 bar. The bubbling at elevated pressure resulted in a higher para-H₂ concentration (ca. 10⁻² M) in the solution.

## 2.3 Quantum chemical calculations

Density functional theory was applied to calculate the J_{BH} coupling constants between the NH and BH protons of BorylCAT–H₂ and ascertain the role of dihydrogen bonding. Coupling constants were calculated for several structures derived via the rotation of the BH₃ group in BorylCAT–H₂. Geometry optimizations and potential energy surface scan calculations were carried out at the oB97X-D/6-311G** level of theory. oB97X-D refers to the dispersion-corrected, range-separated hybrid exchange-correlation functional developed by Chai and Head-Gordon. The NCI (Non-Covalent Interactions) index enables real-space visualization of both attractive and repulsive interactions based on properties of the electron density.

![Fig. 1](image-url)
2.4 Synthesis of ansa-aminoboranes

BorylCAT was synthesized as described previously.\textsuperscript{34} Diaryl-substituted ansa-aminoboranes were prepared as shown in Scheme 2. According to a modified procedure,\textsuperscript{31} treatment of \((2-(2,2,6,6-tetramethylpiperidin-1-yl)phenyl)lithium\) with dimesitylfluoroborane or diphenylchloroborane in diethyl ether gave MesitylCAT and PhenylCAT in 100% and 43% yields, respectively. Addition of boron trichloride to the lithium compound afforded dichloroaminoborane \(A\), which gave iPrPhenylCAT in 94% yield upon reacting with two equivalents of \(\text{o-}^{\text{iso}}\)-isopropylphenyllithium that were generated \textit{in situ} by a Br–Li exchange of \(\text{o-}^{\text{iso}}\)-isopropylbromobenzene.

3. Results

Dihydrogen addition to AABs with different structures, which was investigated in this work using para-\(\text{H}_2\), is represented by Scheme 3.

The scheme does not include nuclear spin dynamics and nuclear spin relaxation in AAB–\(\text{H}_2\) because there is no obvious way to include these processes in chemical equations. Instead, they will be considered in subsequent discussions when it is required.

The \(^{14}\text{N}\)-decoupled \(^1\text{H}\) NMR spectra acquired at 20 °C after bubbling para-\(\text{H}_2\) through the AAB solutions are shown in Fig. 2. In all cases, characteristic antiphase (PASADENA-type) signals were observed for the NH and BH hydrogens in the NMR spectra (signal regions marked with “*” and “#”, respectively). For an AX spin system of spin-1/2 nuclei, a single antiphase doublet would be expected in its spectrum. The shape of the NH signals is close to this conventional shape because the internal multiplet structure introduced by the spin-\(^{14}\text{N}\) nucleus is collapsed due to the self-decoupling caused by a rapid spin relaxation of the \(^{14}\text{N}\) nucleus. The amplitudes of the emissive (e) and absorptive (a) components of the antiphase doublet are not equal. This may be the result of the cross-relaxation occurring during the para-\(\text{H}_2\) bubbling (for \(e > a\))\textsuperscript{28} or low hyperpolarization at low temperatures (for \(e < a\)); see Fig. S1–S4 (ESI\textsuperscript{†}). In contrast, the polarized \(^1\text{H}\) NMR signals of the BH group appear as a set of four antiphase doublets instead of a single doublet because of the splitting induced by the spin-3/2 \(^{11}\text{B}\) nucleus.

In contrast to our previous study,\textsuperscript{28} elevated temperature was not required here to observe polarization effects, meaning that PHIP can be produced at ambient temperatures with the AABs under study. Moreover, these effects were observed even at subzero temperatures (see below, and also Fig. S1–S4 in the ESI\textsuperscript{†}). Comparison of the polarized signal intensities at room temperature reveals that the polarized signal intensity varied in...
the following order: MesitylCat < BorylCat ≈ PhenylCat < iPrPhenylCat. Taking into account the steric hindrance introduced by the substituents at the boryl site, one could draw the conclusion that iPrPhenylCat has optimal properties to produce polarization at room temperature, as the corresponding substituent is located between the highly hindered mesityl group and the moderately hindered phenyl group. However, due to the significant dependence of molecular processes on temperature, the H2 addition process and, consequently, the polarization levels should exhibit a pronounced temperature dependence. To examine this, we performed experiments in which the temperature was varied in the range from −20 °C to 20 °C (see ESI, Fig. S1 to S4). It was found that the intensity of the 1H PHIP effect of the N–H···H–B hydrogen pair decreased gradually with decreasing temperature for BorylCat, iPrPhenylCat and MesitylCat, whereas a clear maximum of PHIP was observed for PhenylCat at ca. 0 °C (Fig. 3).

The existence of the maximum can be explained as the result of an interplay between the nuclear spin relaxation and reaction kinetics, as discussed in detail in Section 4.2.3. Briefly, for low H2 addition rates (slow equilibrium establishment), nuclear polarization is destroyed by the relaxation, whereas at high addition rates, the relaxation is not significant, but the line broadening induced by the exchange diminishes the amplitude of the overlapping antiphase signals. Finally, at very high addition rates (very fast equilibrium establishment), nuclear spin correlation cannot efficiently evolve into observable polarization during the lifetime of the bound state of H2. However, as discussed in Section 4.2.3, this simple explanation cannot exactly explain the position of the maximum if the experimentally measured relaxation times are taken into account.

The reaction kinetic parameters and the relaxation times were measured using the spin saturation transfer (SST) method. The data are shown in Table 1. The details of the determination of the parameters are explained in Section 2.2.

4. Discussion

4.1 Dihydrogen-bond-mediated J-coupling in AABs as a prerequisite for PHIP

In this study, we used the so-called PASADENA experimental scheme, i.e., the reaction with para-H2 was performed in a relatively high magnetic field (7 T). Under these conditions, the weak-coupling regime is fulfilled for all AAB–H2 adducts of para-H2, with the resonance frequency differences between the –NH and –BH protons (>1.2 kHz) in the adducts largely exceeding the J-coupling scalar constants (<6 Hz). The initial nuclear spin correlation of parahydrogen protons, also called the scalar spin order (~1SNH,1BH), transforms upon adduct formation into the longitudinal two-spin order (~1SNH,yBH) because the transverse spin angular momentum components vanish due to decoherence induced by the resonance frequency differences. The 1SNH,1BH two-spin order can be transformed into the two-spin single quantum coherence orders (I1SNH,1yBH,1SNH,1xBH) by a strong radio-frequency pulse of arbitrary phase, and a maximal amplitude is achieved if the pulse flip angle is equal to 45°. The latter orders cannot be observed in NMR directly, as NMR detection is only sensitive to single-quantum single-spin orders (I1SNH,1yBH,1SNH,1xBH). Only these spin orders can contribute to the formation of transverse magnetization detected in practice. However, under the action of J-coupling interactions, the two-spin orders can evolve into observable magnetization.

It is clear from the above brief consideration that a nonzero spin–spin scalar constant between the –NH and –BH hydrogens is required to observe PHIP effects, and the observation of these effects (Fig. 2) serves as evidence of a nonzero coupling constant. However, the nature of this scalar coupling constant has been an open question. The J-coupling can be induced indirectly through the –N–C–C–B– molecular chain of the AABs. However, the typical values expected for a long-range J constant are on the order of 1 Hz, which is much smaller than that observed experimentally (2.9 to 5.6 Hz, see Table 1). Previously, we proposed that J-coupling in the H2 adducts of
AABs is likely direct in nature and may be induced by dihydrogen bonding between the NH and –BH hydrogens. DFT calculations carried out for the BorylCAT–H₂ adduct species provided further information on the origin of the J-coupling. The optimized structure of the BorylCAT–H₂ molecule features a short distance between the NH and BH hydrogens (1.44 Å), and the $J_{HH}$ coupling constant is computed to be 6.2 Hz at this geometry.

Computations also reveal that the BH group of the adduct can easily rotate (or undergo a librational motion) around the B–C bond, as the rotational barrier is predicted to be only 3.8 kcal mol⁻¹ on the potential energy surface. We have therefore calculated the coupling constants along the rotational pathway, i.e. as a function of the C₆-C₆-B–H dihedral angle (see Fig. 4a). The results show a large variation of the $J_{HH}$ values along the rotation. The coupling constant decreases rapidly as the H···H distance increases (it is 2.00 Å in the TS structure), suggesting that the J-coupling is induced by direct H···H interaction rather than via the covalent N–C–C–B– linkage. The statistical averaging of the computed values gives 2.1 Hz for the coupling constant, which is in reasonable agreement with the experimentally observed value (2.9 Hz, see Table 1). Note that for other AABs, this rotation has much higher energy.

Table 1  Selected kinetic parameters, relaxation times and scalar coupling constants for BorylCAT, PhenylCAT, iPrPhenylCAT and MesitylCAT compounds as determined by NMR spectroscopy.

| AAB, $\Delta \delta^a$ (kcal mol⁻¹), $\Delta \delta^a$ (cal mol⁻¹) | T [K] | $k_{dis} \times 10^2$ (s⁻¹) | $K_c \times 10^{-3}$ (M⁻¹) | NH $T_1$ [s] | BH $T_1^* [s]$ | $J_{HH}$ (Hz) | $\Delta \delta^f$ (Hz) |
|-------------------------------------------------------------|-----|-----------------|----------------|-------------|----------------|---------|----------------|
| BorylCAT, 24.2 ± 2.9, 16.3 ± 9.1                            | 273 | 0.1 $^b$        | 0.11 $^c$       | 0.73        | — $^d$          | 2.9     | 4.8            |
| 283                                                        | 0.4 $^b$        | 0.05 $^c$       | 0.88        | — $^d$          | 2.9     | 4.7            |
| 293                                                        | 1.1            | 0.03            | 1.07        | — $^d$          | 2.9     | 4.9            |
| 303                                                        | 7.6            | 0.01            | 1.21        | — $^d$          | —       | —              |
| PhenylCAT, 21.3 ± 0.6, 15.0 ± 2.0                           | 273 | 5.1             | 7.1           | 0.24        | 0.16           | 5.3     | 6.2            |
| 283                                                        | 38.2           | 2.1             | 0.28        | 0.2           | 5.3     | 5.9            |
| 293                                                        | 160.0          | 0.9             | 0.32        | 0.25           | 5.6     | 6.2            |
| 303                                                        | 540.2          | 0.3             | 0.38        | 0.29           | —       | —              |
| iPrPhenylCAT, 22.9 ± 0.9, 13.4 ± 2.8                        | 273 | 0.2 $^b$        | 3.2 $^c$       | 0.23        | 0.15           | 4.8     | 7.6            |
| 283                                                        | 1.0 $^b$        | 1.2 $^c$        | 0.30        | 0.16           | 4.8     | 6.5            |
| 293                                                        | 3.8            | 0.5             | 0.35        | 0.18           | 5.0     | 6.5            |
| 303                                                        | 17.9           | 0.2             | 0.42        | 0.23           | —       | —              |
| MesitylCAT, 23.2 ± 7.7, 16.0 ± 25.7                         | 273 | 0.5 $^b$        | 0.014 $^c$    | 0.34        | 0.16           | 4.3     | 6.2            |
| 283                                                        | 2.3 $^b$        | 0.005           | 0.38        | 0.17           | 4.2     | 4.9            |
| 293                                                        | 10.4           | 0.002           | 0.42        | 0.20           | 4.2     | 4.9            |
| 303                                                        | 32.2           | 0.001           | 0.49        | — $^d$          | —       | —              |

---

4.2 Reaction kinetics, nuclear spin relaxation and PHIP with AABs

In terms of the time scales involved, successful observation of PHIP effects (PASADENA in this work) requires a certain balance between reaction kinetics, nuclear spin relaxation and coherent spin dynamics to convert unobservable nuclear spin orders into observable ones. Any imbalance (e.g., slow or very fast kinetics, slow spin dynamics or fast relaxation) will significantly decrease the observable polarization. The role of coherent spin dynamics requiring a nonzero J-coupling constant was discussed in Section 4.1. There is, however, another requirement: the value of the J-coupling constant must not be much smaller than the linewidth of the multiplet components, or the antiphase components will cancel each other out. The linewidth and the amplitude of the longitudinal two-spin order, in turn, are determined by the kinetic and relaxation parameters. To discuss this complex picture, it is convenient to consider the processes before and after the RF excitation pulse separately.

4.2.1 Evolution before RF pulse. Before the pulse, the experimental protocol included the bubbling of para-H₂ through AAB solutions. This process can be regarded as a stationary process that introduces a flux of para-H₂ nuclear scalar spin order (correlation) into the system. The correlation is transferred further to the AAB adduct with H₂ via intermolecular chemical exchange. A formal kinetic equation for the intermolecular exchange reads

$$\text{AAB} + \text{H}_2 \xrightarrow{k_{\text{ass}}} \frac{k_{\text{ass}}}{k_{\text{dis}}} \text{AAB} - \text{H}_2$$  (1)

where $k_{\text{ass}}$ is the association rate constant and $k_{\text{dis}}$ is the dissociation constant.

As was discussed in Section 4.1, the nuclear spin correlation originating from para-H₂ transforms into the two-spin order...
conditions of para-H₂ bubbling. Moreover, we assume that the para-H₂/AAB/AAB–H₂ system is in chemical and magnetic equilibrium before the bubbling and that the Boltzmann polarization can be ignored due to its low amplitude. Then, if we mark the adducts molecules carrying non-equilibrium nuclear spin orders (e.g. \( I^{NH}_i \), \( I^{BH}_i \)) with "**"", the kinetic equation determining the pumping process will be

\[
\frac{d[AAB - H_2^*]}{dt} = k_{dis} ([AAB - H_2]_0 - [AAB - H_2^*])
\]

\[
- \left( \frac{1}{T^{BH}_1} + \frac{1}{T^{BH}_2} \right) [AAB - H_2^*]
\]

Here, the first term on the right-hand side represents the \( I^{NH}_i, I^{BH}_i \) flux into AAB–H₂, the second term determines the destruction of the order by the nuclear relaxation, \( [AAB-H_2]_0 \) is the total concentration of the adduct, and \([AAB-H_2^*]\) is the instant concentration of the adduct molecules carrying the \( I^{NH}_i, I^{BH}_i \) nuclear spin order. Note that the relaxation process is assumed to be uncorrelated, meaning that the –NH and –BH protons relax independently. If, for simplicity, we group the relaxation constants into \( 1/T^* = 1/T^*_{BH} + 1/T^*_{NH} \), the solution of eqn (2) reads

\[
[AAB - H_2^*](t) = [AAB - H_2^*]_0 \frac{1 - \exp \left(- \frac{1 + k_{dis} T^*_1}{T^*_1} \right)}{1 + \frac{k_{dis} T^*_1}{T^*_1}}
\]

One can see that the stationary state is approached with a function of the sum of the rates of chemical kinetics and nuclear spin relaxation, whereas the amplitude of the achievable \( I_i^{NH}, I_i^{BH} \) nuclear spin order is scaled by the factor determined by the parameter \( k_{dis} T^*_1 \). Note that \([AAB-H_2]_0\) and \( k_{dis} \) (and, much less sensitively, \( T^*_1\)) are temperature dependent. \([AAB-H_2]_0\) is essentially the equilibrium concentration of AAB–H₂ at a given temperature; thus, it can be expressed in terms of the reaction equilibrium constant \( K_e \) through the equation

\[
K_e = \frac{[AAB - H_2]_0}{[AAB-H_2]_0}
\]

After solving [AAB–H₂] from this equation and substituting it in eqn (3), the concentration of molecules carrying the \( I^{NH}_i, I^{BH}_i \) nuclear spin order over the bubbling time is determined by the following expression:

\[
[AAB - H_2^*](t) = [AAB]_0 \frac{1 - \exp \left(- \frac{1 + k_{dis} T^*_1}{T^*_1} \right)}{1 + \frac{1}{k_{dis} T^*_1}}
\]

where \([AAB]_0\) is the total concentration of AAB in both forms (AAB and AAB–H₂). It can be seen that \( K_e \) and \([H_2]_0\) (which is essentially determined by the solubility of H₂ at a given temperature) influence the number of polarized molecules.

It should be noted that eqn (5) assumes that the para-H₂ singlet order is transformed to \( I^{NH}_i, I^{BH}_i \) with 100% efficiency in AAB–H₂ molecules; however, in general, this does not hold. Heuristically, eqn (5) can be equipped with the \( I_i^{NH}, I_i^{BH} \) amplitude
factor \( f_z \) derived in the ESI‡ for chemically exchanging para-H\(_2\) under stationary conditions for a two-spin system. Finally, the following formula for the maximum amplitude of \( \text{I}^\text{NH}_2,\text{I}^\text{BH}_2 \) nuclear spin order acquired in the para-H\(_2\) bubbling process is obtained:

\[
\left[ \text{I}^\text{NH}_2,\text{I}^\text{BH}_2 \right]_{\text{max}} = f_z \cdot \frac{[\text{AAB}]_\Sigma}{1 + \frac{1}{K_z[H_2]_0} + \frac{1}{k_{\text{dis}}T_1^*}} \tag{6}
\]

The amplitude factor is

\[
f_z = \frac{\Delta \omega^2}{\Delta \omega^2 + J^2 + \left( \frac{k_{\text{dis}}}{2\pi} \right)^2} \tag{7}
\]

where \( \Delta \omega \) is the resonance frequency difference between the –NH and –BH protons in Hz, whereas \( J \) is the scalar coupling constant between these protons in Hz. It should be noted that the derivation of eqn (7) assumed that \( k_{\text{dis}} \) is non-zero (see ESI†). The \( f_z \) factor is very close to unity for all AAB adducts considered in this work and can be omitted in the following discussion. In general, however, we note that eqn (6) accounts for the spin dynamics (first term, \( f_z \)), determining the efficiency of the singlet nuclear spin order transformation to \( \text{I}^\text{NH}_2,\text{I}^\text{BH}_2 \); the chemical equilibrium (second term), determining the size of the AAB-H\(_2\) pool; and, finally, the relaxation and the chemical exchange (third term), determining the overall amplitude. After the bubbling is stopped, the nuclear spin system will relax to magnetic equilibrium, reaching the Boltzmann polarization.

4.2.2 Evolution after the RF pulse. We assume that a \( \pi/4 \) RF pulse is applied just after the bubbling is stopped. As discussed in Section 4.1, this pulse will maximize the transverse magnetization that resulted from the INH component at certain temperatures in experiments with BorylCAT (see Fig. 2a and Fig. S1a, b in the ESI† and the discussion below).

4.2.3 Discussion of the experimental results. The experimental data can be analyzed in terms of eqn (11). As the \( f_z \) factor is very close to unity (maximum) for any AAB studied in this work, it can be discarded from consideration. For the other three terms, we can identify the important dimensionless parameters determining the maximum observable antiphase PHIP signal as \( \pi J/(1/T_2^* + k_{\text{dis}}) \), \( K_0[H_2]_0 \) and \( k_{\text{dis}}T_1^* \) for the first, third and fourth terms, respectively. These parameters were calculated and are given in Table 2 for selected temperatures. It should be noted that the total AAB concentration \([\text{AAB}]_\Sigma\) was the same in all cases and thus is not taken into consideration. However, it is obvious that increasing the AAB concentration will proportionally increase the antiphase PHIP signal. Numerical values of the PHIP signal attenuation factors corresponding to the given parameters are presented in Table S1 in the ESI.†

The first parameter, \( \pi J/(1/T_2^* + k_{\text{dis}}) \), determines the strength of the signal self-cancellation due to the line broadening of the overlapping antiphase components. Low values of \( J \) or a large sum of \( 1/T_2^* + k_{\text{dis}} \) will significantly diminish the observable antiphase PHIP signal. For the given AABs, the parameter did not exceed 0.9, meaning that at least 53% (see eqn (11) and Table S1 in the ESI†) of the polarized signal amplitude was lost due to the line broadening effect. The variation of this parameter from one AAB to another is not very significant, and we can conclude that the signal attenuation effect due to the line broadening was on the same order for all AABs. It should be noted that FWHM (\( \Delta \nu \)) values presented in Table 1 for the –NH

\[
\Delta \nu \approx \frac{1}{\pi} \left( \frac{1}{T_2^*} + k_{\text{dis}} \right) \tag{10}
\]

where \( J \) is the scalar coupling constant and \( \Delta \nu \) is the full line width at half maximum (FWHM) in Hz, determined as:

\[
\text{Signal} \propto \frac{1}{1 + \frac{1}{K_z[H_2]_0} + \frac{1}{k_{\text{dis}}T_1^*}} \cdot f_z \cdot \frac{[\text{AAB}]_\Sigma}{1 + \frac{1}{K_z[H_2]_0}} \tag{11}
\]

In addition to the line broadening effect, shape distortions such as accumulation of a complex phase may occur.\(^{47,48}\) These effects, however, were not strong for the –NH and –BH signals. The exception was the signal of ortho-H\(_2\), which had a dispersion component at certain temperatures in experiments with BorylCAT (see Fig. 2a and Fig. S1a, b in the ESI† and the discussion below).
Table 2. Characteristic parameters calculated for the AABs under study and the corresponding NMR signal enhancements observed in the PHIP experiments at selected temperatures

| AAB            | T (K) | $\gamma/(1/T_{1s} + k_{dis})$ | $K_{[H_2]}$ | $k_{dis}T_1s \times 10^4$ | Enh. factor$^a$ |
|----------------|-------|------------------------------|-------------|---------------------------|-----------------|
| BorylCAT       | 273   | 0.6                          | 1.1         | 4.4                       | 1               |
| 283            | 0.6   | 0.5                          | 18.5        | 2                         |                 |
| 293            | 0.6   | 0.3                          | 58.8        | 8.4                       |                 |
| 303            | —     | 0.1                          | 459.8       | —                         |                 |
| PhenylCAT      | 273   | 0.8                          | 71          | 49.0                      | 20              |
| 283            | 0.9   | 21                           | 445.7       | 8                         |                 |
| 293            | 0.9   | 9                            | 2245.6      | 4                         |                 |
| 303            | —     | 3                            | 8885.1      | —                         |                 |
| iPrPhenylCAT   | 273   | 0.6                          | 32          | 2.0                       | 1.2             |
| 283            | 0.7   | 12                           | 10.6        | 5                         |                 |
| 293            | 0.8   | 5                            | 45.2        | 14                        |                 |
| 303            | —     | 2                            | 266.0       | —                         |                 |
| MesitylCAT     | 273   | 0.7                          | 0.14        | 5.7                       | 3               |
| 283            | 0.9   | 0.05                         | 26.1        | 13                        |                 |
| 293            | 0.9   | 0.02                         | 140.9       | 17                        |                 |
| 303            | —     | 0.01                         | 457.3       | —                         |                 |

$^a$ Experimentally measured enhancement factors, i.e., ratios of hyperpolarized (after para-H$_2$ bubbling) and thermally polarized (after relaxation) signals.

The proton resonances are on the order of 8 Hz, meaning that the sum of $1/T_{1s}$ and $k_{dis}$ is ca. 25 Hz (see eqn (10)). This relatively large number implies that the line broadening is primarily determined by $T_{1s}$, nuclear spin relaxation in the AAB molecules because $k_{dis}$ does not exceed 1 Hz for the considered temperatures.

The second parameter, ($K_{[H_2]}$), determines the number of available AAB–H$_2$ molecules. This parameter controls the amplitude of the signal through the chemical equilibrium determined by thermodynamics. Larger $K_{[H_2]}$ values or larger para-H$_2$ concentrations increase the overall signal. The former can be adjusted somewhat by varying the temperature, while the latter can be increased by performing experiments at high para-H$_2$ pressures to increase the concentration. Note that under the given experimental conditions, the parameter $K_{[H_2]}$ varied in the following sequence: PhenylCAT > iPrPhenylCAT > BorylCAT > MesitylCAT (Table 2), meaning that thermodynamically, PhenylCAT is the best candidate for PHIP.

The third parameter ($k_{dis}T_{1s}$) is an intensive property which characterizes the strength of nuclear spin relaxation with respect to the rate of chemical exchange. A low value of this parameter indicates that either the nuclear spin relaxation is very fast or the reaction activity is weak, which leads to significant diminishing of the possible polarization levels. For every AAB in this study, this parameter was not high; however, it was more favorable in the case of PhenylCAT (Table 2).

The above discussion of the experimental data shows that the dimensionless parameters presented in Table 2 conveniently characterize the observed PHIP effects. They also can provide a good strategy basis for maximizing the polarization effects by using eqn (11) if the corresponding kinetic, thermodynamic and NMR parameters are known. Namely, the best results would be expected when $\gamma/(1/T_{1s} + k_{dis}) > 4$, as the arctan function varies weakly once the argument value exceeds 4. Also, other parameters should be $K_{[H_2]} > 1$ and $k_{dis}T_{1s} > 1$ to minimize the respective denominators in eqn (11). These prerequisites, however, are not easy to fulfill in practice. The analysis has shown that the parameters for the considered AABs are far from optimal (see Table 2).

Moreover, the mechanistic features of H$_2$ activation may not be simple to describe. In particular, eqn (11) predicts that the maxima of antiphase PHIP signals should be observed at a certain temperature because $k_{dis}T_{1s}$ decreases significantly at low temperatures, whereas $\gamma/(1/T_{1s} + k_{dis})$ (and possibly $K_{[H_2]}$) decrease significantly at high temperatures. This happens as a result of the exponential dependence of $k_{dis}$ on temperature and qualitatively could explain the distinctive PHIP signal dependence obtained for PhenylCAT (Fig. 3). The position of the PHIP signal maximum, however, cannot be explained quantitatively using eqn (11). Fig. 5 compares the experimentally observed temperature of the PHIP signal amplitude maximum and the theoretically predicted temperature dependence of the PHIP signal amplitude for PhenylCAT. The fact that the positions of the two maxima do not match implies that some hidden dynamics occurs which displaces the experimental maximum to lower temperatures compared to the expected position. The nature of this effect is not clear; however, it may be related to the increase of thermodynamic variance at higher temperatures, defrosting a non-pairwise H$_2$ activation path. The latter process can occur in parallel to the pairwise path; however, in contrast, it does not lead to PHIP effects. Thus, the activation of a non-pairwise path could shift the temperature of the PHIP maximum toward lower temperatures. Another reason for the signal maximum shift may be the temperature sensitivity of the dihydrogen bond in AAB–H$_2$ molecules. The weakening of the dihydrogen bond at higher temperatures can cause the $J$-coupling between the NH and BH protons to vanish, leading to the disappearance of the PHIP maximum.
effect (see Section 4.1). Additionally, in the derivation of eqn (11), we assumed that para-H2 dissolution during the bubbling process is very fast, which may not be a good approximation and may cause some deviations.

Temperature dependence with a PHIP amplitude maximum was observed only for PhenylCAT; for the other AABs, the amplitude of the polarized signals gradually increased with increasing temperature (see Fig. 2 and Fig. S1–S4 in ESI†). This observation is in agreement with the theoretical expectations based on the parameters for BorylCAT, MesitylCAT and iPrPhenylCAT (Table 2). Higher temperatures are required to observe maxima for the latter compounds. BorylCAT, however, is somewhat different from the other AABs. In contrast to the others, at low temperatures, it demonstrated antiphase signals for two unknown polarized species in addition to those for BorylCAT–H2 (see Fig. S1 in ESI†). The attempts to elucidate the structures of these species were not successful because of their very low concentrations. These species could be intermediates, side-products or impurities capable of activating H2 in a pairwise manner, with N and B serving as active FLP centers.

Another peculiar finding for BorylCAT is the observation of a dispersive ortho-H2 signal at elevated temperatures of 283 to 293 K (see Fig. 1a and Fig. S1a–c in the ESI†). This signal looks very similar to a signal with an antiphase shape; however, this shape is impossible for species where two protons have identical resonance frequencies, as in ortho-H2.27 Instead, the observed effect likely arises from a complex phase accumulation induced by a chemical exchange.28 The exchange takes place between the protons captured in BorylCAT–H2 and free H2 molecules. After bubbling para-H2 through BorylCAT solution followed by an RF pulse, free H2 will still be present in the form of NMR-silent para-H2, as it is insensitive to RF pulses. On the other hand, the proton pair in BorylCAT–H2 will be polarized, and the polarization is transferred to free H2 molecules after some time by the chemical exchange. The delayed appearance of ortho-H2 molecules is equivalent to the accumulation of a complex phase for a particular NMR signal, which was observed experimentally in this study. This picture could explain our observations, but only qualitatively.

It was found using numerical spin dynamics simulations that the amplitude of the dispersive signal is very sensitive to the frequency differences between the H2, –NH and –BH resonances (see Fig. S9 in the ESI†), and the large differences present for BorylCAT significantly decrease the amplitude of the dispersive ortho-H2 signal. It is likely that a cross-correlated relaxation in BorylCAT–H2 accompanied by chemical exchange may also have an influence on the dispersive ortho-H2 signal by alternating the signal phase. This was indirectly confirmed by the observation of emissive ortho-H2 signals (signals with negative amplitudes) at almost all temperatures in the experiments with BorylCAT; it can be explained in terms of the cross-correlated relaxation. Evidence that relaxation can convert an originally antiphase polarization to an in-phase (net) polarization is reported in the literature.29,30 Moreover, similar emissive ortho-H2 signals were observed in experiments with other AABs in this work (PhenylCAT), in our previous studies30 on metal-free H2 activation, and also in experiments with diamagnetic Ru complexes.30 As can be seen from the above discussion, a clear understanding of the peculiar subtle effects for ortho-H2 signals requires separate exploration to be considered in full; this will be addressed in future works.

5. Conclusions

In conclusion, we have shown that all AABs studied in this work demonstrate the ability to provide PHIP. Moreover, the required reaction temperatures are close to ambient temperature; thus, the young field of FLP chemistry may bring the efficient production of nuclear spin polarization closer to realization. The experiments are not demanding, and, due to reversibility of H2 addition to the AABs, the polarization can be observed after each cycle of parahydrogen bubbling. The experimentally measured kinetic, thermodynamic and NMR parameters allowed us to analyze the results with the aim of understanding current drawbacks. In particular, it was shown that nuclear spin relaxation accompanied by slow reaction kinetics significantly diminishes the nuclear spin polarization values. Moreover, the rapid nuclear spin relaxation increases the cancellation effect in the antiphase PHIP multiplets. One plausible solution for diminishing the influence of the nuclear spin relaxation would be the replacement of 14N nuclei with 15N. Eliminating the quadrupolar nucleus should increase T1 as well as T2 of the nuclear spins. The reaction kinetics issue can be improved by careful molecular design of the N–B FLP activation centers. The results obtained in this work clarify that variation of the substituents solely at the B-center significantly changes the reaction rates. However, as was found for PhenylCAT, increasing the reaction rates by increasing the temperature may not lead to the expected increase in signal enhancements because the activation of a non-pairwise reaction path or weakening of dihydrogen bonding may occur. Molecular “tuning” at the N-center would provide further flexibility in creating efficient AABs. Thermodynamically, PhenylCAT and iPrPhenylCAT have suitable characteristics, allowing close to quantitative conversion of the AAB molecules to the AAB–H2 adducts, implying high concentrations of the hyperpolarized form of AAB–H2. We believe that further rational investigation of the interactions between para-H2 and AABs can lead to an efficient metal-free para-H2 transfer to various substances and to high NMR signal enhancements. Meanwhile, the mechanistic information provided by PHIP shows that this class of molecules primarily activate H2 molecules in a pairwise manner; demonstrations are already available for six different compounds, including the four studied in this work and two more in a study published earlier.28 We are interested in both generating high levels of hyperpolarization and studying the AABs themselves, as they are unique. The observed hyperpolarization levels are not superior; however, we believe that the relevant improvements described above can provide efficient production of hyperpolarization with AABs. The theoretical analysis presented in the manuscript gives us and other researchers a direction and platform for
further development. This analysis shows that for the systems studied here, the relevant kinetic and NMR parameters are suboptimal for observation of high hyperpolarization; however, this situation can be improved by further design of the AABs.

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