Research Article

Numerical Approaches of the Generalized Time-Fractional Burgers’ Equation with Time-Variable Coefficients

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1. Introduction

The nonlinear convective–diffusive partial differential equations describe various mathematical models in important fields such as heat and mass transfer, fluid mechanics, and engineering. During the last years, various solution methods of ordinary differential equations and partial differential equations have been elaborated.

Burgers’ equation is one of the most important equations involving both nonlinear propagation effects and diffusive effects. A particular form of Burgers’ equation describes the nonlinear wave propagation (the inviscid Burgers’ equation).

Burgers’ equation is a suitable tool for analysis in various fields such as turbulent flows, gas dynamics, shock wave theory, nonlinear wave propagation, longitudinal elastic waves in isotropic solids, sedimentation of polydispersive suspensions and colloids, growth of molecular interfaces, traffic flow, and cosmology [1].

By studying Burgers’ equation with random initial conditions or random forcing, Bec and Khanin [2] explained Burgers’ turbulence. The study of random Lagrangian systems, stochastic partial differential equations, the applications of field theory to the understanding of dissipative anomalies, and of multiscaling in hydrodynamic turbulence are some fields that have significantly benefited from the progress in Burgers’ turbulence. Yu [3] analytically studied the stability and density waves for traffic flow using the perturbation method and shown that the triangular shock waves, soliton wave, and kink wave appear for the density waves.

In the last years, researchers have proved that many phenomena in engineering, bioengineering, physics, and chemistry can be successfully described by mathematical models that use mathematical tools from fractional calculus, i.e., the theory of derivatives and integrals of noninteger order.

Models of viscoelastic materials, Caputo and Mainardi [4]; the signal processing, Marks and Hall [5]; diffusion problems, Olimstead and Handelsman [6]; viscoplastic materials modeling, Diethelm and Freed [7]; mechanical systems subject to damping, Gaul et al. [8]; relaxation and reaction
kinetics of polymers, Glockle and Nonnenmacher [9]; and heat conduction, Hristov [10, 11] are some of the important problems modeled with the help of fractional differential operators.

A very useful collection of numerical algorithms for Caputo-type derivatives, Riemann-Liouville integral operator, and Mittag-Leffler functions is that of Diethelm et al. [12].

In the literature, there are articles in which solutions of Burgers’ equation with different time-fractional derivatives have been determined. We recall some of them. The effects of fractional-order of derivatives on the wave solutions of the generalized Zakharov-Kuznetsov-Burgers equation have been investigated by Faraz et al. [13]. The analytical approximate wave solutions are obtained using the homotopy analysis method and the time-fractional system of Boussinesq-Burgers equations. Using Lie group analysis, the authors derived the infinitesimal groups of transformations, the system of optimal algebras for the symmetry group of transformations, and the similarity variables that reduce the system of fractional partial differential equations to a system of fractional ordinary differential equations. The exact solutions and the physical significance of the solutions are obtained under the invariance condition. Saad et al. [15] have extended the model of the Burgers equation to generalized models based on Liouville-Caputo, Caputo-Fabrizio, and Mittag-Leffler time-fractional derivatives. Using the homotopy analysis transform method, the authors obtained approximate solutions of the newly proposed models.

Balleau and Shiri [16] numerically solved a system of fractional differential equations involving nonsingular Mittag-Leffler kernel using the collocation methods on discontinuous piecewise polynomial space. The existence and regularity of solutions and convergence of the introduced methods are derived.

Recently, Vieru et al. [17] have generalized the time-fractional Atangana-Baleanu derivative. The newly proposed definition contains as particular cases the time-fractional Caputo, Caputo-Fabrizio, and Atangana-Baleanu derivatives.

It is important to note that the operators of fractional derivatives are nonlocal in time and therefore have the advantage of modeling phenomena with memory. Caputo fractional derivatives are nonlocal operators but their kernel is singular. This weakness could have a negative effect when modeling real-world problems.

The fractional derivative operators with the Mittag-Leffler kernel have all the benefits of Caputo operators; in addition, the kernel is nonsingular. Also, their fractional integral operators are the fractional average of the Riemann-Liouville fractional integral of the given function and the function itself. Caputo derivative was conceived for a description of linear short time elastic responses of deformed solids. It was consequently applied to the field of linear viscoelasticity where the Riemann-Liouville derivative was already applied to describe viscoelastic effects. It is known that the asymptotic behaviors of derivative operators with Mittag-Leffler kernel match the power-law behavior. The new fractional operators based on Mittag-Leffler functions have stronger and complex memories allowing capturing behaviors combining simultaneously (crossover) classical diffusion and anomalous behavior. Therefore, to model more complex and nonlinear phenomena, the new operators could be useful tools [11].

In this paper, a nonlinear, one-dimensional, generalized Burgers equation with time-variable coefficients is numerically studied. The generalization consists of considering the fractional differential Burgers’ equation with the generalized time-fractional Atangana-Baleanu fractional derivative with Mittag-Leffler kernel.

A numerical scheme, based on the finite-difference approximations and some integral representations of the two-parameter Mittag-Leffler functions, has been developed along with the consistency, stability, and convergence of the proposed method.

It is important to point out that the studied generalized model can be customized to generate solutions to the problems described by the time-fractional Atangana-Baleanu, Caputo-Fabrizio, and Caputo fractional derivatives.

Numerical solutions of a particular problem with initial and boundary values are determined by employing the proposed method. The numerical results are plotted to compare solutions corresponding to the problems with time-fractional derivatives with different kernels.

2. Preliminary Mathematics

In this section, we present the basic mathematical elements regarding the two-parametric Mittag-Leffler functions and the generalized time-fractional Atangana-Baleanu derivatives. These mathematical notions are necessary for the next sections of this paper.

2.1. One-Parametric and Two-Parametric Mittag-Leffler Functions. The classical one-parametric Mittag-Leffler function is defined as [18, 19]

\[ E_{\alpha}(z) = \sum_{j=0}^{\infty} \frac{z^j}{\Gamma(\alpha j + 1)}, \alpha > 0, \]  

where \( \Gamma(c) = \int_{0}^{\infty} e^{-t} t^{c-1} dt, \Re(c) > 0 \) is the Euler integral of the second kind.

The two-parametric Mittag-Leffler function generalizes the function \( E_{\alpha}(z) \) and is defined by

\[ E_{\alpha,\beta}(z) = \sum_{j=0}^{\infty} \frac{z^j}{\Gamma(\alpha j + \beta)}, \alpha > 0, \beta \in \mathbb{C}. \]  

It is easy to notice that function (1) is a particular case of function (2), so we have

\[ E_{\alpha}(z) = E_{\alpha,1}(z). \]
Let us recall some properties of Mittag-Leffler functions.

\[ E_{\alpha,\beta}(z) = z E_{\alpha,\alpha+\beta}(z) + 1/\Gamma(\beta), \]
\[ E_{\alpha,\beta}(z) = \beta E_{\alpha+1,\beta}(z) + \alpha z \frac{d}{dz} E_{\alpha+1,\beta}(z), \]

(4)

\[ t^\gamma E_{1,\gamma+1}(at) = t^\gamma \sum_{j=0}^{\infty} \frac{(at)^j}{\Gamma(1+j)}, \] - the Miller – Ross function,
\[ t^\gamma E_{1,\gamma+1}(at) = t^\gamma \sum_{j=0}^{\infty} \frac{a^{j(1+\gamma)}}{\Gamma(1+j)} = E_i(y, a) \] - the Robotnov function,

(5)

\[ \int_0^t t^{\alpha-1} E_{\alpha,\beta}(bt^\alpha) dt = z^\beta E_{\alpha,\alpha+1}(bz^\alpha), \beta > 0, \]
\[ \int_0^t t^{\alpha-1}(z-t)^{-1} E_{\alpha,\beta}(bt^\alpha) dt = z^\beta t^{\alpha-1} E_{\alpha,\alpha+1}(bt^\alpha) t(\alpha), \beta > 0, \gamma > 0, \]
\[ \int_0^t (z-t)^{\alpha-1} E_{\alpha,\beta}(bt^\alpha) dt = z^\beta E_{\alpha+1,\beta+1}(bz) \Gamma(\beta), \beta > 0. \]

(7)

The following special form of the one-parametric Mittag-Leffler function [12]:

\[ G(t-r) = E_{\alpha}\left[- \left( \frac{t-r}{y} \right) ^\alpha \right], \alpha \in (0, 1), y > 0, r \in [0, t], \]

(8)

along with its derivative:

\[ M(t-r) = \frac{\partial G(t-r)}{\partial r} = \frac{-1}{t-r} E_{\alpha,0}\left[- \left( \frac{t-r}{y} \right)^\alpha \right], \alpha \in (0, 1), y > 0, r \in [0, t), \]

(9)

has applications in the theory of fractional-order viscoelasticity and in some problems described by fractional differential equations with constant coefficients.

Some numerical algorithms for determining numerical values of the Mittag-Leffler functions have been presented in the reference [12]. These algorithms are based on the integral representations of the Mittag-Leffler functions. We will use in this paper the following integral representations:

If \( \alpha \in (0, 1], \beta \in \mathbb{R}, 0 < \rho < |z|, |\arg z| > \alpha \pi, z \neq 0, \) then

\[ E_{\alpha,\beta}(z) = \int_0^\infty K(\alpha, \beta, x, z) dx + \int_{-\infty}^{\infty} P(\alpha, \beta, \rho, y, z) dy, \]

\[ K(\alpha, \beta, x, z) = \frac{1}{\pi a} x^{\beta-\alpha} \exp \left( -x^{\beta/\alpha} \right) \sin \left( \frac{\pi}{2(1-\beta)} \right) \sin \left( \frac{\pi (1-\beta + \alpha)}{2 \pi x} \right), \]
\[ P(\alpha, \beta, \rho, y, z) = \frac{\rho^{\pi/\alpha}}{2 \pi \alpha} \exp \left( \rho^{\pi/\alpha} \cos \left( \frac{\pi y}{\alpha} \right) \right) \exp \left( i \varphi \right), \]
\[ \varphi = \rho^{\pi/\alpha} \sin \left( \frac{\pi y}{\alpha} \right) + y(1+(1-\beta)/\alpha). \]

(10)

The integral representation

\[ \int_0^\infty e^{-x^\alpha} E_{\alpha,\beta}^{(m)}(\pm \beta x^\alpha) dx = \frac{m! \beta^{\alpha-1}}{(\rho + b)^{\alpha + m}}, \quad \text{Re} (\alpha) > 0, \text{Re} (\beta) > 0, m \in \mathbb{N}, \]

(11)

along with the definition of the Laplace transform of a function \( \mathcal{Y}(t), L\{\mathcal{Y}(t)\} = \int_0^\infty \mathcal{Y}(t) \exp (-st) dt \) give the following relationship:

\[ L\{ t^m E_{\alpha,\beta}^{(m)}(\pm \beta t^\alpha) \} = \frac{m! \beta^{\alpha-1}}{(s^\alpha + b)^{\alpha + m}}, \quad \text{Re} (\alpha) > 0, \text{Re} (\beta) > 0, m \in \mathbb{N}. \]

(12)

In the particular case \( m = 0 \), Equation (12) becomes

\[ L\{ E_{\alpha,\beta}(\pm \beta t^\alpha) \} = \frac{s^{\alpha-1}}{s^\alpha + b}. \]

(13)

2.2 Generalized Atangana-Baleanu Time-Fractional Derivative. The function

\[ \varphi(t, \alpha, \beta) = \frac{1}{1-\alpha} E_{\beta} \left( 1-\alpha \right)^\beta t, t \geq 0, \alpha \in (0, 1), \beta > 0, \]

(14)

is called the generalized Atangana-Baleanu kernel.

The Laplace transform of the kernel (14) is given by

\[ L\{ \varphi(t, \alpha, \beta) \} = \frac{s^{\beta-1}}{(1-\alpha)s^\beta + \alpha}. \]

(15)

Using the Laplace transform, the following properties of the generalized Atangana-Baleanu kernel (14) are found:

\[ L\{ \varphi(t, \alpha, \beta) \} = L\left\{ \lim_{\alpha \to 0} \varphi(t, \alpha, \beta) \right\} = \frac{s^{\beta-1}}{s^\beta} = L\{ 1 \}, \]
\[ L\{ \varphi(t, \alpha, \beta) \} = L\left\{ \lim_{\alpha \to 0} \varphi(t, \alpha, \beta) \right\} = \frac{1}{s^\beta} = L\{ \varphi(0, t) \}, \]
\[ L\{ \varphi(t, \alpha, \beta) \} = L\left\{ \lim_{\beta \to 1} \varphi(t, \alpha, \beta) \right\} = \frac{1}{s^\beta} = L\{ \varphi(0, t) \}, \]
\[ L\{ \varphi(t, \alpha, \beta) \} = L\left\{ \lim_{\beta \to 1} \varphi(t, \alpha, \beta) \right\} = \frac{1}{(1-\alpha)s + \alpha} = L\{ \varphi(0, t) \}, \]

(16)
therefore, 
\[ \lim_{\alpha \to 1} \varphi(t, \alpha, \beta) = \lim_{\beta \to 0} \varphi(t, \alpha, \beta) = 1, \]
\[ \lim_{\alpha \to 1} \varphi(t, \alpha, \beta) = \varphi_1(t, \alpha) = \frac{1}{\Gamma(1-\alpha)} t^{-\alpha}, \]
\[ \lim_{\beta \to 1} \varphi(t, \alpha, \beta) = \varphi_2(t, \alpha) = \frac{1}{1-\alpha} E_\alpha \left( -\frac{\alpha}{1-\alpha} t^{\alpha} \right), \]
\[ \lim_{\alpha \to 1, \beta \to 1} \varphi(t, \alpha, \beta) = \delta(t). \]

In the above relations, functions \( \varphi_0(t, \alpha) \), \( \varphi_1(t, \alpha) \), \( \varphi_2(t, \alpha) \), and \( \delta(t) \) are, respectively, Caputo kernel, Caputo-Fabrizio kernel, Atangana-Baleanu kernel, and the Dirac’s distribution.

**Definition 1.** The generalized Atangana-Baleanu fractional derivative in Caputo sense.

If \( f \in H^1(0, T) \), \( T > 0 \), \( \alpha \in (0, 1] \), \( \beta \in [0, 1] \), the generalized Atangana-Baleanu fractional derivative in Caputo sense, of order \( \alpha \) of the function \( f(t) \), is defined by the relation

\[ \left( GABD_1^{\alpha, \beta} f \right)(t) = \varphi(t, \alpha, \beta) \ast f(t) = \int_0^t \varphi(t - \tau, \alpha, \beta) f(\tau) d\tau. \]

Using Equations (17) and (18), we obtain the following properties of the generalized Atangana-Baleanu time-fractional derivative:

\[ \left( GABD_1^{\alpha, \beta} f \right)(t) = \left( GABD_1^{0, \beta} f \right)(t) = 1 \ast f(t) = \int_0^t f(\tau) d\tau = f(t) - f(0), \]
\[ \left( GABD_1^{1, \beta} f \right)(t) = \delta(t) \ast f(t) = f(t) = \frac{df(t)}{dt}, \]
\[ \left( GABD_1^{\alpha, 1} f \right)(t) = \varphi_1(t, \beta) \ast f(t) = \left( CF D_1^{\beta} f \right)(t), \]
\[ \left( GABD_1^{\alpha, \alpha} f \right)(t) = \varphi_2(t, \beta) \ast f(t) = \left( AB D_1^{\beta} f \right)(t), \]

where \( \left( CF D_1^{\beta} f \right)(t) \) denotes the time-fractional Caputo derivative, \( \left( CF D_1^{\beta} f \right)(t) \) is time-fractional Caputo-Fabrizio derivative, and \( \left( AB D_1^{\beta} f \right)(t) \) denotes the time-fractional Atangana-Baleanu derivative.

Associated with the generalized Atangana-Baleanu derivative, we define the following fractional integral operator:

\[ \left( t_1^{\alpha, \beta} f \right)(t) = (1 - \alpha) f(t) + \alpha \psi_0(t, \beta) \ast f(t), \alpha \in [0, 1], \beta \in (0, 1], \]

where the kernel \( \psi_0(t, \beta) \) is defined as

\[ \psi_0(t, \beta) = \frac{t^{\beta - 1}}{\Gamma(\beta)}. \]

It is observed that \( L\{ \psi_0(t, \beta) \} = 1/s^\beta \), \( \lim_{\beta \to 0} L\{ \psi_0(t, \beta) \} = 1 = L\{ \delta(t) \} \); therefore,

\[ \lim_{\beta \to 0} \psi_0(t, \beta) = \delta(t). \]

Using property (22), the fractional integral operator can be defined for \( \beta = 0 \).

The fractional integral operator (20) has the following properties:

\[ \left( f_1^{1, \beta} f \right)(t) = \delta(t) \ast f(t) = f(t), \]
\[ \left( f_1^{1, 1} f \right)(t) = 1 \ast f(t) = \int_0^t f(\tau) d\tau. \]

Regarding the generalized Atangana-Baleanu derivative and associated fractional integral operator, we remember the following proposition.

**Proposition 2.** The following relationships are fulfilled:

\[ \left( GABD_1^{\alpha, \beta} \left( f_1^{\alpha, \beta} f \right) \right)(t) = f(t) - (1 - \alpha) f(0) \varphi(t, \alpha, \beta), \]
\[ \left( f_1^{\alpha, \beta} \left( GABD_1^{\alpha, \beta} f \right) \right)(t) = f(t) - f(0). \]

The demonstration of the above proposition can be found in the reference [17].

The generalized fractional integral operator (20) contains the following particular cases:

\[ \alpha = 1, \beta \in [0, 1], \]
\[ 1^0 \left( f_1^{1, \beta} f \right)(t) = \psi_0(t, \beta) \ast f(t) = \frac{1}{\Gamma(\beta)} \int_0^t (t - \tau)^{\beta - 1} f(\tau) d\tau, \]

i.e., the well-known Riemann-Liouville fractional integral operator.

\[ \alpha \in [0, 1], \beta = 1, \]
\[ 2^0 \left( f_1^{\alpha, 1} f \right)(t) = (1 - \alpha) f(t) + \alpha \int_0^t f(\tau) d\tau, \]

that is, the integral operator associated to the Caputo-
that is, the fractional integral operator associated with the Atangana-Baleanu fractional derivative.

3. Problem Formulation

The classical one-dimensional Burgers equation with variable coefficients, defined for \((y, t) \in [0, L] \times [0, T], L > 0, T > 0\), is [20]

\[
\frac{\partial u(y, t)}{\partial t} = a(t) \frac{\partial^2 u(y, t)}{\partial y^2} + b(t) u(y, t) \frac{\partial u(y, t)}{\partial y},
\]

where \(a(t) > 0\) and \(b(t)\) are differentiable and bounded functions of the variable \(t\). For \(b(t) = \text{const.}\), Rizun and Engelbrekht [20] have determined the analytical solution of Equation (28).

In the present paper, we consider a generalized form of Equation (28), namely,

\[
G_{AB}D_t^{\alpha, \beta}u(y, t) = a(t) \frac{\partial^2 u(y, t)}{\partial y^2} + b(t) u(y, t) \frac{\partial u(y, t)}{\partial y}, \quad 0 < \alpha, \beta < 1,
\]

where \(G_{AB}D_t^{\alpha, \beta}u(y, t)\) is the generalized Atangana-Baleanu derivative by Equation (18). Along with Equation (29), we consider the initial and boundary conditions

\[
u(y, 0) = \psi(y), \quad y \in [0, L],
\]

\[
u(0, t) = u(L, t) = 0, \quad t \in (0, T].
\]

In the following, we will elaborate a numerical scheme for determining the solution of problem (29)–(31). The proposed scheme is based on the finite-difference method and on the properties of the Mittag-Leffler functions.

Let us consider the discrete set of spatial nodes \(y_k = k\Delta y, k = 0, 1, 2, \ldots, N_1\), respectively, the discrete set of the time nodes \(t_j = j\Delta t, j = 0, 1, 2, \ldots, N_2\), where \(\Delta y = L/N_1, \Delta t = T/N_2\) are the increments of \(y\) and \(t\), respectively.

3.1. Numerical Evaluation of the Generalized Atangana-Baleanu Derivative. Using Equation (18), we have

\[
G_{AB}D_t^{\alpha, \beta}u(y, t) \mid_{(y,t) \to (y_k,t_j)} = \frac{1}{1 - \alpha} \int_0^1 E_{\alpha, \beta}(-\alpha(1 - \alpha)(t_n - t)^\beta) \frac{\partial u(y, t)}{\partial t} dt
\]

\[= \sum_{n=0}^{N_2} \frac{1}{1 - \alpha} \int_0^1 E_{\alpha, \beta}(-\alpha(1 - \alpha)(t_n - t)^\beta) \frac{\partial u(y, t)}{\partial t} dt, \quad n = 1, 2, \ldots, N_2.
\]

The first-order time derivative \(\partial u(y_k, t) / \partial t\), \(t \in [t_j, t_{j+1}]\) is approximated by [21]

\[
\frac{\partial u(y_k, t)}{\partial t} \approx \frac{u(y_k, t_{j+1}) - u(y_k, t_j)}{\Delta t} - \frac{\Delta t^2 \partial^2 u(y_k, t)}{2 \partial t^2} \bigg|_{t=t_j} + O(\Delta t^2).
\]

Using approximation (33), Equation (32) becomes

\[
g_{AB}D_t^{\alpha, \beta}u(y, t) \mid_{(y,t) \to (y_k,t_j)} = \sum_{n=0}^{N_2} \frac{1}{1 - \alpha} \int_0^1 E_{\alpha, \beta}(-\alpha(1 - \alpha)(t_n - t)^\beta) \frac{\partial u(y_k, t)}{\partial t} dt + R_{kn}
\]

\[- = \sum_{n=0}^{N_2} \frac{1}{1 - \alpha} \int_0^1 E_{\alpha, \beta}(-\alpha(1 - \alpha)(t_n - t)^\beta) \frac{\partial u(y_k, t)}{\partial t} dt + R_{kn}
\]

\[
= \sum_{n=0}^{N_2} A_n \nu_k^{n+1} - \sum_{n=0}^{N_2} B_n \nu_k^n + R_{kn},
\]

where

\[
A_n = \frac{(t_{n+1} - t_j)E_{\alpha, \beta}(\lambda(t_n - t_j)^\beta)}{(1 - \alpha)\Delta t}, \quad B_n = \frac{(t_{n+1} - t_j)E_{\alpha, \beta}(\lambda(t_n - t_j)^\beta)}{(1 - \alpha)\Delta t}, \quad \nu_k^n = \nu_k(t_j).
\]

Introducing notations

\[
D_{y_j} = \begin{cases} (n - 1)E_{\alpha, \beta} - nE_{\alpha, \beta}, & j = 0, \\ (n - j)E_{\alpha, \beta} - (n - j)E_{\alpha, \beta}, & j \in (1, 2, \ldots, n - 1), \\ E_{\alpha, \beta}, & j = n, \end{cases}
\]

\[
D_{y_{j+1}} = E_{\alpha, \beta},
\]

where \(E_{\alpha, \beta} = E_{\alpha, \beta}(\lambda(t_n - t_j)^\beta)\), Equation (35) can be written in the equivalent form

\[
G_{AB}D_t^{\alpha, \beta}u(y, t) \mid_{(y,t) \to (y_k,t_j)} = \frac{1}{1 - \alpha} \sum_{n=0}^{N_2} D_{y_j} \nu_k^n + R_{kn}.
\]
The truncation error $R_{kn}$ is defined as

$$R_{kn} = \sum_{j=0}^{n-1} \frac{1}{1-a} \int_{t_j}^{t_{j+1}} \frac{\Delta t^2 \partial^2 u(y_k, \tau)}{\partial \tau^2} \bigg|_{\tau=t_j} E_{\beta} \left( -\frac{\alpha}{1-a} \right) (t_n - \tau)^{\beta} d\tau$$

$$= \frac{\Delta t^2}{2(1-a)} \sum_{j=0}^{n-1} \frac{\partial^2 u(y_k, \tau)}{\partial \tau^2} \bigg|_{\tau=t_j} \left[ (n-j)E_{\beta} - (n-j-1)E_{\beta+1} \right]$$

$$\leq \frac{C_t \Delta t^2}{2(1-a)} \left\{ \max_{0 \leq k \leq n} \frac{\partial^2 u(y_k, \tau)}{\partial \tau^2} \right\}, \tag{39}$$

where $C_t$ is a constant coming from the bounded of the Mittag-Leffler function. The inequality (39) ensures the consistency of the proposed method because, assuming that the function $u(y, t)$ is sufficiently smooth on its domain of definition, the truncation error $R_{kn}$ tends to zero if the time step $\Delta t$ tends to zero. So, at point $(y_k, t_n)$, the generalized Atangana-Baleanu time-fractional derivative is approximated by

$$GAB D_{\beta}^{\alpha, \beta} u(y_k, t) \bigg|_{(y_k, t_n)} = \frac{1}{1-a} \sum_{j=0}^{n} D_{\beta}^{\alpha, \beta} \left( (y_j, t_{j+1}) \right). \tag{40}$$

In the above relations, the numerical values of the Mittag-Leffler functions are evaluated using the integral representation (10).

An important property of the coefficients $D_{\beta}^{\alpha, \beta}$ is given in the following.

**Lemma 3.** Coefficients $D_{\beta}^{\alpha, \beta}$ given by (37) have negative values for $j = 0, 1, \ldots, n - 1$. For $a \in (0, 1/2)$, $D_{\beta}^{\alpha, \beta} \rightarrow 0$ and $D_{\beta}^{\alpha, \beta} \rightarrow 1$ if $\Delta t \rightarrow 0$.

**Proof.** Let $Q(z)$ a function defined as

$$Q(z) = \sum_{j=0}^{n} \frac{\alpha}{1-a} \int_{t_j}^{t_{j+1}} E_{\beta} \left( -\frac{\alpha}{1-a} \right) (t_n - \tau)^{\beta} d\tau = \frac{1}{1-a} \sum_{j=0}^{n} E_{\beta} \left( -\frac{\alpha}{1-a} \right) (t_n - \tau)^{\beta}.$$

(42)

Using the definition of Mittag-Leffler function, $Q(z)$ can be written as

$$Q(z) = \sum_{k=0}^{\infty} \frac{\alpha}{1-a} \int_{t_j}^{t_{j+1}} E_{\beta} \left( -\frac{\alpha}{1-a} \right) (t_n - \tau)^{\beta}.$$

(43)

Using the analytical expression, it is known that Mittag-Leffler function $E_{\beta}^{\alpha, \beta}(z)$ is an increasing function. Because, in this study, parameter $\lambda < 0$, it results that function $Q(z)$ is a decreasing function. The coefficients $D_{\beta}^{\alpha, \beta}$ are written as

$$D_{\beta}^{\alpha, \beta} = [(n-j+1)E_{\beta-j} - (n-j)E_{\beta}] - [(n-j)E_{\beta-j} - (n-j-1)E_{\beta+1}] \rightarrow [(n-j+1)E_{\beta} \lambda - (n-j)E_{\beta} \lambda] - [(n-j)E_{\beta} \lambda - (n-j-1)E_{\beta+1} \lambda].$$

(44)

Since, the function $Q(z)$ is decreasing we obtain that $D_{\beta}^{\alpha, \beta} < 0$, $j = 1, 2, \ldots, n - 1$.

$$D_{\beta}^{\alpha, \beta} = E_{\beta}^{\alpha, \beta} \lambda \Delta t^{\beta}. \quad D_{\beta}^{\alpha, \beta} < 0$$

(45)

Table 1 gives the values of this derivative determined with the analytical expression (45), respectively, with the numerical formula (40) for $t_k = 0.02k, k = 1, 2, \ldots, 20$ and $\alpha = 0.45, \beta = 0.75$.

**Example 2.** Find the solution of the fractional equation

$$GAB D_{\beta}^{\alpha, \beta} u(t) = t^\gamma, \gamma > 0, \alpha, \beta \in (0, 1]. \tag{46}$$

Using the Laplace transform, it is found that the analytical solution of Equation (46) is given by

$$u(t) = (1-\alpha)t^\gamma + \frac{\alpha\Gamma(1+\gamma)}{\Gamma(1+\beta+\gamma)} t^{\beta+1}. \tag{47}$$

(47)

Table 2 gives the values of the solution of Equation (46) determined with the analytical expression (47), respectively, with the numerical formula (40) for $t_k = 0.01k, k = 4, 8, \ldots, 80, \alpha = 0.45, \beta = 0.75, \lambda = 2$. 
3.2. Particular Cases. Note that the expression (40) of the time-fractional derivative remain valid, obviously with the corresponding particularizations of the parameters $a$ and $b$.

In the third case, for $a = 1$, there is an indeterminacy because $\lim_{a \to 1} (a/(1-a)) = \infty$.

To eliminate this indeterminacy, we use the following asymptotic expansion of the Mittag-Leffler function [22]:

$$
\beta \in (0, 1), \quad |z| \to 0, 1, \quad |\arg(z)| \leq \pi,
$$

$$
E_{\beta}(z) = -\sum_{n=1}^{\infty} \frac{1}{\Gamma(1-\beta n)} + O\left(\frac{1}{2^{\beta n} n!}\right), \quad \beta \in (0, 1), \quad n \in \mathbb{N}, \quad n_0 > 1.
$$

Using Equation (44) into (34), we obtain the following relation:

$$
\lim_{\alpha \to 1} \frac{1}{1-\alpha} \int_{t_j}^{t_{j+1}} E_{\beta} \left(\frac{-a}{1-\alpha}(t_n - \tau)^\beta\right) d\tau = \int_{t_j}^{t_{j+1}} \lim_{\alpha \to 1} \frac{1}{1-\alpha} \left(1 - \alpha^\beta(1-\beta)^\beta\right) \left[\Gamma(1-\beta n)\right] d\tau
$$

$$
= \frac{-(t_n - \tau)^{-1-\beta}}{(1-\beta)(1-\beta)} \left| t_{j+1} - t_j \right| + \frac{1}{T(2-\beta)} \left[ (t_n - t_j)^{-1-\beta} - (t_n - t_{j+1})^{-1-\beta} \right].
$$

(49)

As expected, the integrand in (45) is the Caputo kernel $q_{\beta}(t, \beta)$. Replacing (45) into (34), we obtain the following approximate formula for the time-fractional Caputo derivative:

$$
C_{D_t}^{\beta} u(y, t) \big|_{y = y_{j+1}} = \sum_{j=0}^{\infty} C_{n_j} u_j^\beta,
$$

(50)
Table 2: Comparison between analytical and numerical results of the solution of Equation (46).

| $t_k = k\Delta t$ | Analytical solution | Numerical solution | $t_k = k\Delta t$ | Analytical solution | Numerical solution |
|-------------------|---------------------|--------------------|-------------------|---------------------|--------------------|
| $\Delta t = 0.01$ |                     |                    | $\Delta t = 0.01$ |                     |                    |
| 0.04              | 0.0912002 × 10^{-4} | 9.0985187 × 10^{-4} | 0.44              | 0.1277624           | 0.1277683          |
| 0.08              | 3.7158954 × 10^{-3} | 3.7171929 × 10^{-3} | 0.48              | 0.1537559           | 0.1537622          |
| 0.12              | 8.5174136 × 10^{-3} | 8.5192433 × 10^{-3} | 0.52              | 0.1824127           | 0.1824195          |
| 0.16              | 0.0153980           | 0.0154002          | 0.56              | 0.2137889           | 0.2137962          |
| 0.20              | 0.0244342           | 0.0244371          | 0.60              | 0.2479393           | 0.2479471          |
| 0.24              | 0.0356989           | 0.0357023          | 0.64              | 0.2849178           | 0.2849261          |
| 0.28              | 0.0492606           | 0.0492645          | 0.68              | 0.3247773           | 0.3247861          |
| 0.32              | 0.0651852           | 0.0651896          | 0.72              | 0.3675701           | 0.3675794          |
| 0.36              | 0.0835363           | 0.0835411          | 0.76              | 0.4133474           | 0.4133572          |
| 0.40              | 0.1043754           | 0.1043807          | 0.80              | 0.4621601           | 0.4621703          |

where

$$C_n = \frac{1}{\Delta t^\beta} \left[ (n+1)^{-\beta} - n^{-\beta} \right], \quad j = 0,$$

$$C_n = \frac{1}{\Delta t^\beta} \left[ (n+1)^{-\beta} - 2(n)\Delta \beta + (n+j-1)^{-\beta} \right], \quad j = 1, 2, \ldots, n-1,$$

$$C_n = \frac{1}{\Delta t^\beta} \Delta \beta^n.$$

(51)

3.3. Numerical Solution to the Burgers Equation with Generalized Atangana-Baleanu Time-Fractional Derivative.

To determine the numerical solution of Equation (29), we approximate the first- and second-order derivative with respect to $y$ by [23–25].

$$\frac{\partial u(y, t)}{\partial y} \bigg|_{(y,t)} = \frac{u_{k+1}^i - u_{k-1}^i}{2\Delta y} , k = 1, 2, \ldots, N_1 - 1,$$

$$\frac{\partial^2 u(y, t)}{\partial y^2} \bigg|_{(y,t)} = \frac{u_{k+1}^i - 2u_k^i + u_{k-1}^i}{\Delta y^2} , k = 1, 2, \ldots, N_1 - 1.$$

(52)

The nonlinear term $u(y, t) \frac{\partial u(y, t)}{\partial y}$ is replaced by the equivalent term

$$\frac{1}{3} \left[ u(y, t) \frac{\partial u(y, t)}{\partial y} + \frac{\partial u^2(y, t)}{\partial y} \right].$$

(53)

For simplicity of calculations and proving the stability of the scheme, the nonlinear term is approximated by [30, 31]

$$u(y, t) \frac{\partial u(y, t)}{\partial y} \bigg|_{(y,t)} = \frac{1}{3} \left[ u_{i}^{j+1} u_{i+1}^{j+1} - u_{i}^{j+1} u_{i}^{j+1} + u_{i}^{j+1} u_{i+1}^{j+1} - u_{i}^{j+1} u_{i}^{j+1} \right].$$

(54)

Using Equations (40), (52), and (54), we obtain the following numerical scheme for fractional Burgers equation.
The initial and boundary conditions (30) and (31) are transformed in following discrete relationships:

\[u_k^0 = u(y_k, 0) = u(y_k, t_0) = \psi(y_k) = \psi_k, \quad k = 0, 1, \ldots, N_1,\] (56)

\[u_j^0 = u(0, t_j) = u(y_0, t_j) = 0, \quad j = 1, 2, \ldots, N_2,\] (57)

\[u_{N_2}^j = u(L, t_j) = u(y_{N_1}, t_j) = 0, \quad j = 1, 2, \ldots, N_2.\] (58)

Using (57) and (58), the numerical scheme (55) is written in the following metrical form:

\[M(n)U^n = P(n), \quad n = 1, 2, \ldots, N_2,\] (59)

where the matrix \(M(n) = (m_{ij}^n)_{i,j=1,2,\ldots,N_1-1}\) is the following tridiagonal matrix:

\[
M(n) = \begin{pmatrix}
  m_{11}^n & m_{12}^n & 0 & 0 & 0 & \cdots & 0 \\
  m_{12}^n & m_{22}^n & m_{23}^n & 0 & 0 & \cdots & 0 \\
  0 & m_{23}^n & m_{33}^n & m_{34}^n & 0 & \cdots & 0 \\
  \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\
  0 & 0 & \cdots & m_{N_1-2,N_1-3}^n & m_{N_1-2,N_1-2}^n & m_{N_1-2,N_1-1}^n & 0 \\
  0 & 0 & \cdots & 0 & m_{N_1-1,N_1-2}^n & m_{N_1-1,N_1-1}^n & 0
\end{pmatrix}_{(N_1-1,N_1-1)}.
\] (60)

with

\[
m_{ij}^n = \begin{cases}
  b(t_n\delta y) \left( u_i^{n-1} + u_{i+1}^{n-1} \right) + \frac{a(t_n)}{\Delta y^2} \delta_{ij} + \frac{D_{mn}}{1 - \alpha} \frac{2a(t_n)}{\Delta y^2} \delta_{ij}, \\
  b(t_n\delta y) \left( u_i^{n-1} + u_{i-1}^{n-1} \right) + \frac{a(t_n)}{\Delta y^2} \delta_{ij}, \\
  \frac{b(t_n\delta y)}{6\Delta y} \left( u_i^{n-1} + u_j^{n-1} \right) + \frac{a(t_n)}{\Delta y^2} \delta_{ij}, \\
  b(t_n\delta y) \left( u_i^{n-1} + u_{i-1}^{n-1} \right) + \frac{a(t_n)}{\Delta y^2} \delta_{ij}, \\
  \frac{b(t_n\delta y)}{6\Delta y} \left( u_i^{n-1} + u_j^{n-1} \right) + \frac{a(t_n)}{\Delta y^2} \delta_{ij}, \\
  \frac{b(t_n\delta y)}{6\Delta y} \left( u_i^{n-1} + u_j^{n-1} \right) + \frac{a(t_n)}{\Delta y^2} \delta_{ij},
\end{cases}
\]

\[U^n = \left( u_1^n, u_2^n, \ldots, u_{N_1-2}^n, u_{N_1-1}^n \right)^T,\] (61)

\[P(n) = \left( p_1^n p_2^n \cdots p_{N_1-2}^n p_{N_1-1}^n \right)^T,\] (62)
Table 3: Comparison between numerical solutions corresponding to different fractional derivatives with the nonsingular kernel.

| Values of time | (GAB)-fractional derivative for α = 0.1, β = 0.5 | (AB)-fractional derivative for α = 0.1, β = 0.1 | (CF)-fractional derivative for α = 0.1, β = 1.0 |
|----------------|---------------------------------------------|---------------------------------------------|---------------------------------------------|
| 0.10           | 0.022000                                   | 0.021000                                   | 0.022000                                   |
| 0.20           | 0.019000                                   | 0.018000                                   | 0.019000                                   |
| 0.30           | 0.015000                                   | 0.014000                                   | 0.015000                                   |
| 0.40           | 0.012000                                   | 0.011000                                   | 0.012000                                   |
| 0.50           | 0.012000                                   | 0.011000                                   | 0.012000                                   |
| 0.60           | 0.007455                                   | 0.007055                                   | 0.007612                                   |
| 0.70           | 0.006107                                   | 0.005780                                   | 0.006235                                   |
| 0.80           | 0.005106                                   | 0.004833                                   | 0.005213                                   |
| 0.90           | 0.004355                                   | 0.004122                                   | 0.005929                                   |
| 1.00           | 0.003784                                   | 0.003582                                   | 0.004913                                   |

Table 5: Comparison between numerical solutions corresponding to different fractional derivatives with the nonsingular kernel.

| Values of time | (GAB)-fractional derivative for α = 0.4, β = 0.5 | (AB)-fractional derivative for α = 0.4, β = 0.4 | (CF)-fractional derivative for α = 0.4, β = 1.0 |
|----------------|---------------------------------------------|---------------------------------------------|---------------------------------------------|
| 0.10           | 0.029000                                   | 0.028000                                   | 0.033000                                   |
| 0.20           | 0.025000                                   | 0.024000                                   | 0.055000                                   |
| 0.30           | 0.020000                                   | 0.019000                                   | 0.022000                                   |
| 0.40           | 0.016000                                   | 0.015000                                   | 0.017000                                   |
| 0.50           | 0.016000                                   | 0.015000                                   | 0.017000                                   |
| 0.60           | 0.009886                                   | 0.009433                                   | 0.011000                                   |
| 0.70           | 0.008096                                   | 0.007726                                   | 0.009024                                   |
| 0.80           | 0.006768                                   | 0.006460                                   | 0.007555                                   |
| 0.90           | 0.005773                                   | 0.005509                                   | 0.006451                                   |
| 1.00           | 0.005016                                   | 0.004787                                   | 0.005610                                   |

Let us introduce the notations:

\[(u^n_k)_{y} = \frac{u^n_{k+1} - u^n_{k}}{\Delta y}, \quad (u^n_{k})_{\tau} = \frac{u^n_{k+1} - u^n_{k}}{\Delta \tau}, \quad \gamma_1 = \frac{u^n_{k+1} - 2u^n_{k} + u^n_{k-1}}{\Delta y}, \]

\[U^n = \{u^n_k, k = 0, 1, \ldots, N_1\}, \quad (U^n)_{y} = \{ \gamma_1 \}, \quad \|U^n\| = \left\| \sum_{k=1}^{N_1-1} u^n_{k} v^n_{k}, \|U^n\| = \sum_{k=1}^{N_1-1} u^n_{k}^2, \|U^n\|_{\infty} = \lim_{k \to \infty} |u^n_k| \right\}.

The following boundedness theorem will be proved.

**Theorem 6.** Assume that \(\psi \in H^2_{0}[0, 1], \alpha \in (0, 1/2), \beta \in [0, 1].\) There exists a constant \(K\) such that the numerical solution derived by the finite difference scheme (55) satisfies inequality \(\|U^n\| \leq K, \quad n = 1, 2, \ldots, N_2.\)

**Proof.** Equation (55) can be written in the equivalent form

\[D^n_{\varepsilon} u^n_{k} + \frac{1}{\alpha} \sum_{j=1}^{N_1} D^n_{\alpha} u^n_{j} + \frac{D^n_{\beta} u^n_{k-1}}{\alpha} - \frac{b(t^n_{k}) \left[u^n_{k-1} - (u^n_{k} - u^n_{k-1}) + u^n_{k} - u^n_{k-1}\right]}{6\Delta y} = \frac{u^n_{k-1} - 2u^n_{k} + u^n_{k+1}}{\Delta y^2}.

Multiplying Equation (65) by \(\Delta y u^n_{k}\) and summing for \(k = 1\)
Equation (71) implies that

$$
\frac{D_{mn}}{2(1-\alpha)} \left( \|U^n\|^2 - \|U^0\|^2 \right) \leq -a(t_n) \|U^n\|^2 \leq 0,
$$

(72)

therefore,

$$
\|U^n\| \leq \|U^0\|.
$$

(73)

But, there exist a constant $\bar{K}$ such that $\|U^0\| \leq \bar{K}$, respectively. $\|U^n\| \leq \bar{K}$. Theorem 6 shows that for a given function $\psi(y) = u(y,0)$, the numerical solution of the problem is bounded as time increases.

In order to prove a stability result, we consider a perturbed problem, i.e., the fractional equation (29) with a different initial condition $\psi(y) = u(y,0)$.

We say that the numerical method is **globally stable** if there exists a constant $\bar{K}$ such that $\|U^n - U^n\| \leq \bar{K} \|\psi(y) - \psi(y)\|$, where $U^n$ and $U^n$ are solutions corresponding to initial conditions $\psi(y)$ and $\psi(y)$, respectively. Using the boundedness theorem one can prove that the numerical scheme (55) corresponding to the Equation (29) is globally stable.

**Theorem 7.** (Convergence). Assume that $\{u^n_1\}$, $\{\tilde{u}^n_1\}$ are the solutions of the difference scheme (55) and (29) and $\{\epsilon^n_k\} = \tilde{u}^n_k - u^n_k$. There exists a positive constant $C$ such that $\|E^n\| \leq C(\Delta t + \Delta y^2)$.

**Proof.** Using the equalities (55) and (63), we obtain

$$
\frac{D_{mn}}{1-\alpha} \epsilon^n_k^2 - a(t_n) \epsilon^n_k^2 = \frac{1}{1-\alpha} \sum_{j=1}^{n} (-D_{n_j}) \epsilon^n_k^2 + \frac{\epsilon^n_k}{1-\alpha} \left( \frac{\Delta y^2}{2} \sum_{j=1}^{n} (-D_{n_j}) \epsilon^n_k^2 \right)
$$

$$
\leq \frac{1}{2(1-\alpha)} \sum_{j=1}^{n} (-D_{n_j}) \left( \|U^n\|^2 + \|U^0\|^2 \right) + \frac{(-D_{n_j})}{2(1-\alpha)} \|U^n\|^2
$$

$$
= \frac{1}{2(1-\alpha)} \sum_{j=1}^{n} (-D_{n_j}) \left( \|U^n\|^2 + \|U^0\|^2 \right)
$$

$$
+ \frac{1}{2(1-\alpha)} \left( \sum_{j=1}^{n} (-D_{n_j}) \right) \|U^n\|^2
$$

$$
= \frac{1}{2(1-\alpha)} \sum_{j=1}^{n} (-D_{n_j}) \|U^n\|^2 + \frac{(-D_{n_j})}{2(1-\alpha)} \|U^n\|^2 + \frac{D_{mn}}{2(1-\alpha)} \|U^n\|^2.
$$

(70)

Using (70) and mathematical induction, we obtain that $\|U^n\|^2 \leq \|U^0\|^2$ for $j = 1, 2, \ldots, n - 1$. Therefore, we have

$$
\frac{D_{mn}}{1-\alpha} \|U^n\|^2 - a(t_n) \|U^n\|^2 \leq \frac{1}{2(1-\alpha)} \sum_{j=1}^{n} (-D_{n_j}) \|U^n\|^2 + \frac{(-D_{n_j})}{2(1-\alpha)} \|U^n\|^2
$$

$$
\leq \frac{1}{2(1-\alpha)} \sum_{j=1}^{n} (-D_{n_j}) \|U^n\|^2 + \frac{(-D_{n_j})}{2(1-\alpha)} \|U^n\|^2 = \frac{D_{mn}}{2(1-\alpha)} \|U^n\|^2.
$$

(71)

$\square$
Using Equation (69) and the Cauchy-Schwarz inequality 
\[ e_k^2 e_k^2 \leq (1/2)(e_k^2)^2 + (e_k^2)^2, \]
we get
\[
\begin{align*}
\frac{D_m}{1 - \alpha} \| E^n \|^2 + a(t_n) \| (E^n) \|^2 &\geq \frac{1}{2(1 - \alpha)} \sum_{k=1}^{n} (-D_m) \| E^n \|^2 + \frac{1}{2(1 - \alpha)} \left( \sum_{k=1}^{n} (-D_m) \right) \| E^n \|^2 \\
&\quad + (-D_m) \| E^n \|^2 + (T^n, E^n) + (S^n, E^n) + (R^n, E^n).
\end{align*}
\]
(77)

We note that \( \| E^n \| = 0 \), and \( \sum_{j=0}^{n-1} (-D_{nj}) = D_m \). Based on the Cauchy-Schwarz inequality, we obtain following inequalities:
\[
(R^n, E^n) \leq \gamma_1 e_n \| R^n \|^2 + \frac{1}{4 \gamma_1 e_n} \| E^n \|^2,
\]
\[
(T^n, E^n) \leq \frac{b(t_n)}{3} \sum_{k=1}^{n} \left[ \left( u^k_k - u^k_k \right) u^k_k - \left( u^k_k \right) u^k_k \right] e_k
\]
\[
= \frac{b(t_n)}{3} \sum_{k=1}^{n} \left[ \left( u^k_k - u^k_k \right) u^k_k - \left( u^k_k \right) u^k_k \right] e_k
\]
\[
= \frac{b(t_n)}{3} \sum_{k=1}^{n} \left[ \left( u^k_k - u^k_k \right) u^k_k - \left( u^k_k \right) u^k_k \right] e_k
\]
\[
\leq \frac{b(t_n)}{3} \left[ \frac{1}{\| E^n \|^2} \| E^n \|^2 \right] + \epsilon_n \| (E^n) \|^2 + \frac{1}{4 \gamma_1 e_n} \| E^n \|^2.
\]
(78)

Taking \( e_n = 3a(t_n)/4b(t_n) \delta_n \), \( \gamma_1 = 1/3, \gamma_2 = \gamma_3 = 1/2 \), we obtain the following inequality:
\[
\delta_n \| E^n \|^2 + a(t_n) \| (E^n) \|^2 \leq \frac{1}{2(1 - \alpha)} \sum_{j=1}^{n-1} (-D_m) \| E^n \|^2 + \pi_n \| E^{n-1} \|^2 + \pi_n \| E^{n-1} \|^2
\]
\[
+ k_1 \sigma \| E^n \|^2,
\]
(79)

where
\[
\rho_n = \frac{K b(t_n)}{6 a(t_n)} \left[ 6 + 2K b(t_n) + a(t_n) \right],
\]
\[
\pi_n = \frac{K b(t_n)}{18 a(t_n)} \left[ 4K b(t_n) + 3a(t_n) \right],
\]
\[
\delta_n = \frac{D_m}{2(1 - \alpha)} - \rho_n.
\]

Now, we will show that
\[
\delta_n \| E^n \|^2 + a(t_n) \| (E^n) \|^2 \leq \frac{k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))},
\]
(81)

by mathematical induction. For \( n = 1 \), inequality (81) holds according to (79). Suppose that
\[
\delta_n \| E^n \|^2 + a(t_n) \| (E^n) \|^2 \leq \frac{k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))}, j = 1, 2, \ldots, n - 1,
\]
(82)

then,
\[
\| E^n \|^2 \leq \frac{k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))}, j = 1, 2, \ldots, n - 1.
\]
(83)

Introducing (83) into (79), we obtain
\[
\delta_n \| E^n \|^2 + a(t_n) \| (E^n) \|^2 \leq \frac{1}{2(1 - \alpha)} \sum_{j=1}^{n-1} (-D_m) \| E^n \|^2 + \pi_n \| E^{n-1} \|^2
\]
\[
+ \frac{k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))}
\]
\[
+ \pi_n \| E^{n-1} \|^2 + \frac{k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))} + \frac{k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))}
\]
\[
+ k_1 a(t_n) \delta_n \| \Delta t + \Delta y \|^2
\]
(84)

Therefore, it results in
\[
\| E^n \|^2 \leq \frac{k_1 a(t_n) \| \Delta t + \Delta y \|^2}{\delta_n - \pi_n - (D_{nn} + D_{nn})/(2(1 - \alpha))},
\]
(85)

Using the discrete Sobolev’s inequality \( \| E^n \|_\infty \leq C_1 \| E^n \| + C_2 \| (E^n) \|, \) we obtain the result of Theorem 7.

Some comparisons between our results and other results from the literature are presented at the end of this section.

(a) If we consider \( a(t) = 1, b(t) = -1, \alpha = 1, \beta \in (0, 1), \) then the problem given by (29)–(31) is identical to the problem studied by Li et al. [23], for \( p = 1 \)

(b) If \( a(t) = V, b(t) = -1, \alpha = \beta, \) our solutions are identical to that obtained by Yadav et al. [24] for \( g(x, t) = 0 \)

(c) If \( \alpha = \beta, \) the numerical scheme (38) for the generalized Atangana-Baleanu derivative is identical with that given in [25], Equation (3.2).
4. Example

In this section, we investigate the problem described by the fractional differential equation (29), with \( a(t) = 0.01 + 0.1 \sin^2(\pi t/4) \), \( b(t) = -10 \sin(\pi t/4) \), along with the following initial and boundary conditions:

\[
\begin{align*}
u(y, 0) &= \psi(y) = 2(1 - \exp(-y)), y \in [0, 1], \\
u(0, t) &= u(0.5, t) = 0.
\end{align*}
\]

Equation (86)

The computational results are taken at the time step size \( \Delta t = 1.0 \). The initial and boundary conditions: are plotted in Figure 4 for different types of fractional derivatives with the non-singular kernels. Because the kernels of fractional derivatives play an essential role in describing the effects of memory, Figure 3 shows the profiles of the kernels corresponding to the generalized Atangana-Baleanu, Atangana-Baleanu, and Caputo-Fabrizio derivatives, versus fractional parameter \( \alpha \), for \( t = 0.3 \). It is observed that for \( \alpha < 0.8 \), the values of the generalized kernel (GAB) are higher than the (AB) ones but smaller than the (CF) ones; therefore, the different damping effects will be transferred on the values of the function \( u(y, t) \).

The numerical solutions of Burgers equation with generalized Atangana-Baleanu time-fractional derivatives (GAB) are plotted in Figure 4 for different fractional parameters \( \alpha \) and \( \beta \), for \( y \in [0, 0.5] \), and for different values of the time \( t \). The computational results are taken at the time step size \( \Delta t = 10^{-2} \) and spatial step size \( \Delta y = 10^{-2} \).

Figure 4 shows that, for the considered cases, the function \( u(y, t) \) is decreasing in relation to time and it is increasing in relation to the fractional parameters \( \alpha \). The curves in Figure 4 indicate that the numerical solutions are bounded, and the numerical method is stable.

A comparison between the solutions of the Burgers equation with the generalized Atangana-Baleanu derivative, the Atangana-Baleanu derivative and the Caputo-Fabrizio derivative is presented in Tables 3–5, for \( x = 0.4 \), and for different values of time \( t \).

As expected, due to the properties of the three nonsingular kernels, the values of the solution corresponding to the Atangana-Baleanu derivative are the lowest, respectively, and the values of the solution corresponding to the Caputo-Fabrizio derivative are the highest.

5. Conclusion

A numerical scheme to solve the fractional Burgers equation with variable coefficients has been developed.

The proposed algorithm is a linear implicit finite difference scheme.

Unconditional stability and convergence of the method are proved for the fractional parameters \( \alpha \in (0.1/2), \beta \in (0, 1) \).

The proposed scheme is suitable to solve fractional Burgers’ equation with four fractional derivatives, namely, generalized Atangana-Baleanu, Atangana-Baleanu, Caputo-Fabrizio, and Caputo derivatives.

Data Availability

There is no any data available.

Conflicts of Interest

The authors declare that they have no conflicts of interest.

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