Relationship between large-scale ionospheric field-aligned currents and electron/ion precipitations: DMSP observations

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Abstract
In this study, we have derived field-aligned currents (FACs) from magnetometers onboard the Defense Meteorological Satellite Project (DMSP) satellites. The magnetic latitude versus local time distribution of FACs from DMSP shows comparable dependences with previous findings on the intensity and orientation of interplanetary magnetic field (IMF) Bző and Bẑ components, which confirms the reliability of DMSP FAC data set. With simultaneous measurements of precipitating particles from DMSP, we further investigate the relation between large-scale FACs and precipitating particles. Our result shows that precipitation electron and ion fluxes both increase in magnitude and extend to lower latitude for enhanced southward IMF Bẑ, which is similar to the behavior of FACs. Under weak northward and southward Bẑ conditions, the locations of the R2 current maxima, at both dusk and dawn sides and in both hemispheres, are found to be close to the maxima of the particle energy fluxes; while for the same IMF conditions, R1 currents are displaced further to the respective particle flux peaks. Largest displacement (about 3.5°) is found between the downward R1 current and ion flux peak at the dawn side. Our results suggest that there exists systematic differences in locations of electron/ion precipitation and large-scale upward/downward FACs. As outlined by the statistical mean of these two parameters, the FAC peaks enclose the particle energy flux peaks in an auroral band at both dusk and dawn sides. Our comparisons also found that particle precipitation at dawn and dusk and in both hemispheres maximizes near the mean R2 current peaks. The particle precipitation flux maxima closer to the R1 current peaks are lower in magnitude. This is opposite to the known feature that R1 currents are on average stronger than R2 currents.

Keywords: Field-aligned currents, Aurora, Particle precipitation, DMSP

Introduction
Auroral field-aligned currents (FACs), also known as Birkeland currents, are an important transport mechanism for energy and momentum between the magnetosphere and ionosphere and thus are of fundamental importance for understanding the solar wind–magnetosphere–ionosphere–thermosphere coupling (e.g., Milan et al. 2017). The distribution of FACs largely depends on the orientation of the interplanetary magnetic field (IMF), while the intensity is also modulated by the solar wind electric field and ram pressure, as well as ionospheric conductivity (Korth et al. 2010). Solar irradiation flux variations can slightly affect the intensity of FACs, e.g., through modulation of ionospheric conductivity, but the solar wind forcing has been demonstrated to dominate FAC distributions (e.g., Edwards et al. 2017). For southward IMF conditions, the FACs exhibit a consistent pattern of two concentric rings flowing in and out of the ionosphere, labeled as poleward Region 1 (R1) and equatorward Region 2 (R2) currents (Iijima and Potemra 1976a, b), while for northward IMF condition, an additional poleward current system is often observed in the dayside cusp region, known as “NBZ” or “Region 0” currents (e.g., Iijima et al. 1984; Milan et al. 2000).
The global distribution of FACs and their variations have been widely investigated using magnetometer measurements on board low Earth orbital (LEO) satellites, like Triad (e.g., Armstrong and Zmuda 1975; Iijima and Potemra 1976a,b), OGO 5 (Orbiting Geophysical Observatory) (e.g., Sugiuira 1975), Freja (e.g., Lühr et al. 1996), Dynamics Explorer 2 (e.g., Weimer 2001), Ørsted (e.g., Papitashvili et al. 2002), or CHAMP (Challenging Minisatellite Payload) (e.g., Wang et al. 2005). However, as pointed out by Lühr et al. (1996), computing FACs along track from such single-satellite mission has always suffered from non-uniqueness. As the satellite moves through three-dimensional regions of high current density, the recorded magnetic field changes can be interpreted in terms of current density only if certain assumptions are made: the currents are static over the time of spacecraft measurements and have a simple geometric configuration, e.g., the sheet of FACs is elongated approximately east–west, and perpendicular to the satellite orbit (Lühr et al. 2020). Later with the development of satellite constellations, like Cluster II (e.g., Escoubet et al. 2001), ST5 (e.g., Slavin et al. 2008; Gjerloev et al. 2011), as well as the recent Swarm mission (e.g., Ritter et al. 2013; Dunlop et al. 2015; Trenchi et al. 2019), FACs’ density estimations are made using simultaneous measurements from multi-point data spanning a two-dimensional area in space. By comparing FACs derived from single- and dual-satellite of Swarm, Lühr et al. (2015) found that at the auroral latitudes, the large-scale FAC signatures are consistent between the two approaches, but discrepancies commonly appear poleward of 75° magnetic latitude (MLAT). Filamentary FAC signatures are captured by the dual-satellite approach at poleward auroral oval under northward IMF condition, but are missed by the single-satellite technique. Another application of a multi-satellite approach to calculate FACs is the Active Magnetosphere and Planetary Electrodynamics Response Experiment (AMPERE) project (e.g., Anderson et al. 2000; Green et al. 2009; Korth et al. 2010). Two-dimensional average configurations of FACs are derived at cadences of less than an hour from magnetic measurements of 66 Iridium satellites.

FACs can be carried by either electrons or ions, but the dominant carriers are thought to be magnetospheric/ionospheric electrons flowing along the magnetic field lines. It was observed that contribution to FACs from downward precipitating electrons exceeds that of upward flowing ionospheric ions typically by an order of magnitude (e.g., Cattell et al. 1979). Several studies have been carried out for resolving the relationship between FACs and particle precipitations. Focusing on large-scale upward FACs and electron precipitations, Korth et al. (2014) have compared simultaneous observations of FACs derived from AMPERE and N₂ Lyman–Birge–Hopfield (LBH) auroral emission measured by the Global Ultraviolet Imager (GUVI) on board the Thermosphere, Ionosphere, and Mesosphere Energetics and Dynamics (TIMED) satellite. The electron precipitation is inferred from a statistical relation between LBH intensity and electron energy flux (Sotirelis et al. 2013). They found that for southward IMF, the electron precipitation occurred primarily within and near large-scale upward currents, while such correspondence between upward currents and electron precipitation is less evidenced for northward IMF condition. In another study, Robinson et al. (2018) have reexamined the relation between FACs from AMPERE and LBH auroral emission from GUVI, but used a different empirical function as that from Zhang and Paxton (2008) to convert the LBH intensity to electron energy flux. They have combined the upward and downward currents, and found that the particle energy flux correlates well with FACs intensity, showing strongest dependence near magnetic midnight hours. Carter et al. (2016) have investigated FACs from AMPERE with the aurora ultraviolet (UV) emission from Imager for Magnetopause-to-Auroral Global Exploration (IMAGE) mission. They found that the R2 current is more closely aligned with the distribution of auroral UV emission than the R1 current, whether that be in the discrete auroral zone of dusk or in the post-midnight diffuse aurora sector. This finding is different from the result of Korth et al. (2014), who concluded that the upward R1 current at dusk has a better dependence on the electron precipitation than the upward R2 current at dawn. Therefore, the relation between FACs and particle precipitations is still an open issue, which encourages further investigation.

In the studies of Korth et al. (2014) and Carter et al. (2016), they have compared FACs and particle precipitations derived from different missions, the AMPERE and GUVI as well as AMPERE and IMAGE. When comparing measurements between different missions, the separations in temporal and spatial coverage between the mission pair might lead to differences in interpreting their results. For example, Carter et al. (2016) tried to make their comparison consistent by sorting the FACs and IMAGE emissions into bins under similar IMF orientations, as there is no overlap between the AMPERE and IMAGE missions. Therefore, applying direct comparisons between FACs and particle precipitations from the same platform allows for temporal overlap. This study utilizes simultaneous magnetic and particle precipitation measurements from the Defense Meteorological Satellite Project (DMSP). The magnetic measurements of DMSP have earlier been used for investigating, e.g., the dayside FAC source regions (Wing et al. 2010), magnetic perturbations at polar cap during geomagnetic storms
(Knipp et al. 2014), and pointing flux in the dayside polar cap boundary regions (Lu et al. 2018). Here, we derive statistical analysis of FAC signatures from DMSP, especially from the recent F16–F18 satellites. Additionally, the DMSP data set allows investigating contributions from precipitating electron and ions separately. In “Dataset Section”, we shortly introduce the DMSP mission as well as its magnetic and precipitation measurements. In “Observations Section”, we first verify the reliability of FAC signatures derived from DMSP by comparing to nearly simultaneous Swarm measurements. Then, the IMF dependence and relation between the FACs and precipitation particles observed by DMSP are discussed. Evaluations against previous studies are provided in “Discussion Section”. Finally, we summarize the main findings in “Summary Section”.

Dataset
The special sensor magnetometer (SSM) and special sensor J (SSJ) measurements from DMSP as well as the magnetic measurements from Swarm

DMSP is a series of Earth-observing spacecraft since 1962. In this study, we used observations from satellites F-16, F-17, and F-18, that fly in Sun-synchronous, near-polar orbits, with inclinations of about 98.8° and periods of about 102 min (Burke et al. 2011). The three DMSP satellites all carry the SSM, which is a triaxial fluxgate magnetometer consisting of three separate cores to measure the magnetic field components in a range of ±63,553 nT and with an accuracy of 2 nT (Kilcommons et al. 2017). Alken et al. (2014) have calibrated the SSM dataset from DMSP F16–F18 for modeling the Earth’s main field. The calibrated dataset have been extended until the end of 2016 by Alken et al. (2020). The SSM data are provided with time resolution of 1 Hz. These data are used for deriving FACs in this study. Measurements from the Special Sensor J (SSJ/5) instrument sense precipitating electrons and ions with energy flux between 30 eV and 30 keV. SSJ/5 data from DMSP F16–F18 satellites during a 5-year period (2010–2014) have been calibrated by Redmon et al. (2017) and are publicly accessible at https://satdat.ngdc.noaa.gov/dmsp/data/f18/ssi/j/.

Figure 1 (middle panel) shows the altitude evolution of DMSP F-16 (black), F-17 (blue), and F-18 (red) from 2009 to the end of 2016. For this period, Alken et al. (2020) provided calibrated magnetic data from DMSP. The three satellites fly at about 850 km, and their local time coverages (bottom panel) are in the evening for ascending nodes (thick-solid lines) and in the morning for descending nodes (thin-dotted lines). Their orbits drift very slightly in local time. In “Justification of FACs derived from DMSP Section”, we show one example of the Swarm observations to compare with DMSP data, and thus, the altitude and local time evolution of Swarm A has also been added, here marked in green. Swarm A flies at about 460 km altitude which is slowly decaying. The orbit of Swarm A drifts in local time, which needs about 133 days for covering 24-h local times (see also Xiong et al. 2016). On Swarm, the magnetic field is measured by a vector field magnetometer (VFM) and an absolute scalar magnetometer (ASM), and the latter one is used to calibrate the vector field. The VFM is a high precision instrument capable of magnetic field measurements with time resolution up to 50 Hz and a precision of about 0.1 nT. In the example, we use magnetic time series of Swarm with time resolution of 1 Hz.

\textbf{FACs derived from DMSP magnetic measurements}

To isolate ionospheric sources in geomagnetic measurements, geomagnetic field contributions from the Earth’s core, crust, and magnetosphere are estimated by the CHAOS-6 model (Finlay et al. 2015) and removed from the data. The satellite-based residual observations are applied to estimate radial and field-aligned currents in the ionosphere. The basic equation for deriving electric currents from magnetic measurements is Ampère’s law and it is for the vertical current component $j_z$:

$$j_z = \frac{1}{\mu_0} \left( \frac{db_y}{dx} - \frac{db_x}{dy} \right)$$  \hspace{1cm} (1)

where $b_x$ and $b_y$ are the magnetic field components caused by the currents, $x$ and $y$ are the coordinates in northward and eastward direction of the NEC frame (north, east, center), respectively, and $\mu_0$ is the permeability of free space. The radial currents are then transferred to FACs by projection into the direction of the local magnetic field line (see also Ritter et al. 2013).

The residual magnetic field data from DMSP have been applied in a processor originally developed for the Swarm mission, to derive FAC data set of DMSP. Since the temporal resolutions of the Swarm and DMSP time series are both 1 Hz, no additional adaptation was applied in the processor. Stolle et al. (submitted to Earth Planets and Space) reported that large-scale FACs are well captured by non-high-precision magnetometer in space. The DMSP FAC data of F16–F18 during 2009–2016 are available at ftp://magftp.gfz-potsdam.de/DMSP/FAC/. For verification with DMSP FACs, we apply an example of Swarm as provided as daily product through ESA.

Particle precipitation data from DMSP F-16 to F-18 provided by Redmon et al. (2017) covers a 5-year period from 2010 to 2014 (indicated by light-orange shadow in Fig. 1). Therefore, we focused also on this 5-year period to investigate DMSP FAC data. The solar flux variation for this period is shown in the top panel of Fig. 1. The
mean value of the P10.7 index between 2010 and 2014 is 116.3 sfu.

**Observations**

**Example of magnetic and particle precipitations data measured by DMSP**

Figure 2 shows a high-latitude orbital segment of DMSP F-18 around 16:45 UTC on 5 February 2014. The first panel shows the three components of magnetic residuals to the CHAOS-6 model predictions, $\Delta b_x$, $\Delta b_y$, and $\Delta b_z$ pointing to the north, east, and center of the Earth. Both $\Delta b_x$ and $\Delta b_y$ show fluctuations of about $\pm$ 200 nT above 70º magnetic Apex latitude (Emmert et al. 2010, hereafter called MLAT), while $\Delta b_z$ is much smoother. Occasional step-like jumps in $\Delta b_z$ are remainders from the calibration process (Alken et al. 2020); however, they do not affect the FACs calculation [see Eq. (1)]. As expected, FACs of amplitudes within $\pm 5$ $\mu$A/m² in the second panel are derived where $\Delta b_x$ and $\Delta b_y$ show small-scale fluctuations.

The SSJ/5 instrument on board DMSP uses 20 channels detecting precipitating electrons and ions. Differential energy fluxes derived from each channel are added up to get the integrated energy flux. It is then divided by the number of observed counts to get an average energy of electrons and ions (Hardy et al. 2008; Redmon et al. 2010). The lower panels of Fig. 2 show the integrated electron energy flux, the electron average energy (AVG), the integrated ion energy flux, and the ion average energy. An elevated level of precipitating electrons appears at around $\pm$ 70º MLAT and stays higher inside the polar cap than equatorward of the precipitation maxima. The ion precipitation data are sparser; however, an increased level is still recognized at a few degrees higher latitude than for the electron precipitation in this example. Here, the precipitating electrons and ions fluxes reach to lower latitudes than the intense FACs. In the study, we compare FAC pattern with the electron and ion energy fluxes, but not with the average energy. Carter et al. (2016) and Korth et al. (2014) similarly compared derived electron energy fluxes with FAC patterns.
Justification of FACs derived from DMSP

Figure 3 shows FAC observations along the same DMSP F-18 orbital segments on 5 February 2014 together with the Swarm A orbit that has been spatially and temporally close. Both satellites have been counter-rotating at similar magnetic local times, and the UTC difference at the highest magnetic latitude of their respective orbit was about 14 min. Enhanced FAC signatures display at
similar magnetic latitudes. The 1 Hz FAC time series by Swarm A including small-scale features shows slightly larger amplitudes (panel b). However, after applying a low-pass filter with a cut-off frequency of 20 s to the 1 Hz data (panel c) from both satellites, the large-scale structures show consistent features with similar amplitudes. Outside the auroral region, the noise level of the current intensity derived from Swarm A is within ±0.02 μA/m², while the DMSP-derived currents have a noise level of 0.15 μA/m². Considering the largest amplitude of FAC on the dusk side of this event is about ±3 μA/m², this corresponds to 0.7 and 5% for Swarm A and DMSP, respectively. However, after applying a low-pass filter with cut-off frequency of 20 s, the levels of intrinsic fluctuations derived from both satellites are less than 0.02 μA/m². This example suggests that FACs derived from DMSP compare well with observations from high precision data and are quite reliable. Similar conclusions can be derived from comparisons with DMSP F-16 and F-17 (not shown here).

Subsequently, we display the IMF dependence of FACs from DMSP. The data from all three DMSP satellites have been combined to achieve a wider local time coverage. A low-pass filter with a cut-off frequency of 20 s was applied to the 1 Hz FAC time series of DMSP to reflect large-scale FAC signatures. The data were sorted into MLAT (1º) and MLT (1 h) bins and plots for different levels of IMF magnitudes and orientations were created. Figure 4a shows the MLAT versus MLT distributions of the averaged FAC signatures from DMSP in the northern hemisphere. The subpanels are arranged by the intensity of IMF $B_y$ and $B_z$ components: from left to right, $B_y$ component varies from 10 to -10 nT, and from top to bottom, $B_z$ varies from 10 to -10 nT, both in steps of 5 nT. Due to the low occurrence of data for $B_y$ or $B_z$ beyond ±10 nT, results are not shown here. The lack of data around noon and midnight hours are due to DMSP’s dawn–dusk orbits. For each subpanel, dusk/dawn is on the left/right, and noon/night is on the top/bottom. For enhanced southward $B_y$, the intensity of FACs increases, the R1 and R2 currents appear in a clearer pair pattern, and they expand to lower latitudes. For northward $B_y$, the known current pair NBZ appears poleward of the R1 sheet around local noon. Concerning the dependence on $B_y$, for example, the upward R1 current tends to extend from afternoon to noon hours when $B_y$ is more negative. Figure 4b shows the distribution of FACs in the southern hemisphere, which basically shows similar dependence on the IMF $B_z$. The NBZ current is not monitored, as there is almost no data coverage around noon hours in the southern hemisphere. However, coverage is enhanced around midnight in the southern hemisphere. The IMF dependences of FAC shown in Fig. 4a, b compare well to those of previous publications (e.g., Wang et al. 2008; Korth et al. 2010; Milan et al. 2017; He et al. 2012; Coxon et al. 2014; Carter et al. 2016; Laundal et al. 2018). Furthermore, the FACs intensity in the northern hemisphere is slightly higher than that in the southern hemisphere, which is consistent with the findings of Coxon et al. (2016) and Workayehu et al. (2019). Another feature seen here is that the intensity of R1 currents is stronger than for R2 currents at both dawn and
Fig. 4  a. The MLAT versus MLT distribution of FACs derived from the DMSP satellites for the northern hemisphere. The subpanels are arranged by the intensity of IMF $B_y$ and $B_z$ components: from left to right, the $B_y$ component varies from 10 to $-10$ nT, and from top to bottom, the $B_z$ varies from 10 to $-10$ nT, both with a step of 5 nT. b. As the same as Fig. 4a, but for the FACs distribution in the southern hemisphere.
Fig. 4 (continued)
Fig. 5  a) The MLAT versus MLT distribution of electron energy flux derived from the DMSP satellites for the northern hemisphere. b) As the same as Fig. 5a, but for the electron energy flux distribution in the southern hemisphere.
Fig. 5 (continued)
dusk sides, which is consistent with the previous findings (e.g., Christiansen et al. 2002; Le et al. 2010; Wang et al. 2019). All these comparisons further support the reliability of the DMSP FACs.

Comparison between FACs and particle precipitations observed by DMSP

Figure 5a, b shows the distribution of electron energy flux from DMSP in the two hemispheres, respectively. Similar to FACs, the electron energy flux increases with stronger southern \( B_z \) and expands to lower latitude. During northward \( B_z \), the electron energy flux is much weaker at dusk than at dawn, and this feature is true for both hemispheres. Figure 6a, b shows the distribution of ion energy flux from DMSP in the two hemispheres, respectively. It also increases with southward \( B_z \). However, its intensity is weaker at dawn than at dusk, being an opposite dusk–dusk asymmetry to the electron energy flux.

In the following analysis, we will concentrate on the morning and evening sector. Figure S1 of the appendix shows that data coverage of DMSP is uniform in these sectors. Figure 7 shows latitudinal profiles of DMSP FACs (black), and energy flux of electrons (blue) and ions (red) along the dusk–dawn meridian and averaged within \( \pm 1.5 \) h MLT. The solid and dashed lines represent the values in the northern and southern hemispheres, respectively. The left and right columns show results for intense northward (5~10 nT) and southward (\(- 10 \sim -5 \) nT) \( B_z \), respectively, and with varying \( B_z \) condition for each subpanel. For visualization purpose, the values of electron and ion precipitations have been divided by factors of \( 1 \times 10^{12} \) and \( 1 \times 10^{11} \) eV/cm\(^2\)/ster/s, respectively. The intensity of all three quantities increase and they cover a more equatorward latitudinal range for southward \( B_z \) than for northward \( B_z \). For northward \( B_z > 5 \) nT, however, it is not straightforward to reliably separate the R1/R2 current along the dusk–dawn meridian, though the peaks of precipitating electrons and ions are well discernable. Therefore, the following comparison between FACs and precipitation will be conducted for \( B_z < 5 \) nT. The intensity of the ion flux is in average stronger on the dusk side, while the intensity of the electron energy flux is stronger in the dawn side. The peaks of electron energy flux at the dusk and dawn sides are located around \( 70^\circ \) MLAT and \( 64^\circ \) MLAT, while the peaks of the ion energy flux are located around \( 65^\circ \) and \( 68^\circ \) MLAT, respectively. In the following, peaks of all three quantities are identified, as is demonstrated in Fig. 8 for the profiles of lower right panel in Fig. 7. As expected, the locations of the peaks of the upward FACs correspond well to the peaks of electron energy flux, and the locations of the peaks of the downward currents correspond well to the peaks of the ion energy flux. Table 1 shows the locations (in MLAT) of peaks of upward/downward FACs and electron/ion energy flux for all subpanels with \( B_z < 5 \) nT, separately for the dusk and dawn sides as well as two hemispheres.

Table 2 shows the mean difference of all entries in Table 1, separately for the dusk and dawn sides and the two hemispheres. Positive values reflect that the electron/ion energy flux peak is located poleward of the upward/downward FACs peak, and negative values reflect its equatorward location. On average, larger differences occur between the locations of R1 currents and ion/electron flux peaks (especially at the dawn side), than for R2 currents. At R1, the particle flux peaks are poleward of the FAC peaks, while, at R2, the particle flux peaks are equatorward of the FAC peaks. Our results are in several points consistent with the findings of Carter et al. (2016), but in some points different from that of Korth et al. (2014). This will be discussed in more detail in “Discussion Section”. Furthermore, Fig. 9 visualizes the mean locations of these peaks. The red/blue thick lines represent the upward/downward FACs, and the dot/cross within black circles represent the precipitation electron/ions, respectively. Interestingly, the FAC peaks enclose the particle energy flux peaks at both dusk and dawn side and in both hemispheres.

Discussion

Earlier observations showed that there is generally an imbalance between the pair of opposite-flowing R1 and R2 FACs in both dawn and dusk sides, i.e., the total current flowing in R1 FACs exceeds that in R2 FACs, and as a result, there are net currents flowing into (out of) the ionosphere at the dawn (dusk) sides (e.g., Le et al. 2010). This feature is also evidenced in the DMSP measurements, as shown in Fig. 4. Previous theories predicted that R1 currents come from the low-latitude boundary layer of closed field lines or the high-latitude boundary layer of open field lines (Sonnerup, 1980; Tanaka, 1995), while the R2 currents originate in the inner magnetosphere or at the inner edge of the plasma sheet (Southwood, & Harel et al. 1981). Through this picture, the R1 currents are more directly linked to the IMF and respond in general faster to IMF variations than the R2 currents by tens of minutes (Kikuchi et al. 2000). As also seen here, the large-scale R1 and R2 currents appear much clearer as an opposite-flowing pair and their intensity increase for southward \( B_z \). It also means that the absolute magnitude of an R1/R2 imbalance is larger for stronger southward \( B_z \) and contributes to a significant part of the IMF-averaged imbalance between R1 and R2 currents. The intensity of precipitating electron and ion flux are shown in Figs. 5, 6, 7 for the same IMF conditions. Here, the electron energy flux is stronger at dawn (corresponding to the dawn-side
Fig. 6  a The MLAT versus MLT distribution of ion energy flux derived from the DMSP satellites for the northern hemisphere.  b As the same as Fig. 6a, but for the ion energy flux distribution in the southern hemisphere.
Fig. 6 (continued)
Fig. 7 The latitudinal profiles of FACs (black), energy flux of electron (blue) and ion (red), and quantities are averaged within ±1.5 h MLT centered on the dusk–dawn meridian. The solid and dashed lines represent the values in the northern and southern hemispheres, respectively. Here, we show only the results for the more northward (left, $B_z = 5–10$ nT) and more southward (right, $B_z = -10--5$ nT) $B_z$ conditions.
upward R2 currents) than at dusk (corresponding to the dusk-side R1 currents), while the ion energy is stronger at dusk (corresponding to the dusk-side downward R2 currents) than at dawn (corresponding to the dawn-side downward R1 currents). This is in contrast to R1 currents being stronger than R2 currents. Although this observation is evident from the data, we can currently not provide a good explanation on the mechanism. However, it suggests that the R1/R2 currents are not solely determined by the intensity of downward particle energy flux.

Korth et al. (2014) found that under southward $B_z$, the distribution of electron energy flux agrees well with the dusk-side upward R1 current, but, on the dawn side, the upward R2 current shows wider latitudinal coverage and the peak is located a few degrees equatorward than that of the electron energy flux. They explained that, at dusk side, the electron precipitation is mainly due to field-aligned acceleration, while it is not the case for the dawn side.

**Table 1** The MLAT of peaks for upward/downward FAC, electron energy flux, and ion energy flux, separately for the dusk and dawn sides, as well as for the two hemisphere. Only the conditions with $B_z < 5$ nT have been considered.

| IMF   | Dusk | Dawn |
|-------|------|------|
|       | Bz (nT) | By (nT) | FAC (R1) | Electron flux up | Ion flux up | FAC (R2) | Electron flux down | Ion flux down |
| NH    | 0–5   | 5–10  | 74 67 74 68 | 66 75 67 70 |
|       | 0–5   | 5–10  | 76 69 74 69 | 66 76 68 71 |
|       | – 5–0 | 78 68 73 70 | 68 76 68 71 |
|       | – 10–5 | 74 67 72 68 | 68 73 68 69 |
|       | 5–10  | 73 66 71 67 | 66 73 66 70 |
|       | 0–5   | 74 67 73 68 | 67 73 67 70 |
|       | – 5–0 | 74 67 72 68 | 68 74 68 71 |
|       | – 10–5 | 72 65 71 67 | 66 74 67 70 |
| SH    | 0–5   | 5–10  | 72 62 69 63 | 63 69 64 69 |
|       | 0–5   | 5–10  | 71 64 69 65 | 64 72 66 69 |
|       | – 5–0 | 70 63 68 63 | 64 72 66 69 |
|       | – 10–5 | 70 62 68 64 | 64 71 64 68 |
|       | 5–10  | 73 66 72 68 | 68 74 67 70 |
|       | 0–5   | 75 68 73 69 | 69 75 69 72 |
|       | – 5–0 | 75 68 73 69 | 69 76 69 71 |
|       | – 10–5 | 73 66 72 68 | 67 74 67 70 |
|       | 5–10  | 73 65 71 67 | 67 74 66 69 |
|       | 0–5   | 74 67 72 68 | 68 74 68 70 |
|       | – 5–0 | 74 67 72 68 | 68 74 68 70 |
|       | – 10–5 | 72 65 71 66 | 66 73 66 69 |
|       | 5–10  | 70 63 67 64 | 64 73 65 69 |
|       | 0–5   | 70 63 69 64 | 65 71 65 68 |
|       | – 5–0 | 70 63 69 64 | 65 71 65 68 |
|       | – 10–5 | 70 62 70 62 | 64 70 65 68 |
precipitating electrons in the post-midnight to dawn sector where large population of diffused aurora electrons is observed (Meredith et al. 2003; Newell et al. 2009). Therefore, Korth et al. (2014) suggested that the distribution of upward current in dawn sector does not well agree with the distribution of precipitating electrons. However, this conclusion does not fully agree with our findings. As listed out in Table 2, on average, the electron energy flux shows smaller displacement with respect to the upward R2 current at dawn, compared to the dusk upward R1 current. Another difference is that, on both dusk and dawn sides, DMSP observed precipitating electrons show wider latitudinal extension than the upward FACs, while in Fig. 6 of Korth et al. (2014), the dawn-side upward R2 currents have a longer tail extending to lower latitudes.

Table 2 The mean differences between the peaks of electron/ion energy flux and upward/downward FACs separately for dusk and dawn sides and two hemispheres. Only the conditions with Bz < 5 nT have been considered.

|          | Dusk                                           | Dawn                                           |
|----------|------------------------------------------------|------------------------------------------------|
|          | Electron flux and upward FAC (R1) | Ion flux and downward FAC (R2) | Electron flux and upward FAC (R2) | Ion flux and downward FAC (R1) |
| NH       | $-2.0^\circ$ | $1.1^\circ$ | $0.6^\circ$ | $-3.4^\circ$ |
| SH       | $-1.8^\circ$ | $1.1^\circ$ | $0.2^\circ$ | $-3.8^\circ$ |

Such difference might come from the FACs data processing from AMPERE. Carter et al. (2016) explained that the spatial resolution afforded by the spherical harmonic analysis of the historical Iridium data over a period of 2 h may not resolve the current structures measuring less than 3° in latitude. In their analysis of Iridium data with 10 min resolution, they showed that the latitude extent of the currents is generally narrower than those obtained from 2 h resolution. The DMSP FACs are derived from 1 Hz magnetic data and the FACs have then be represented by a 20 s cut-off filter. The direct comparison with electron precipitation shows similar equatorial boundaries at dawn (see Figs. 7 and 8) and a high agreement of the location of the peaks of the R2 upward currents and electron precipitation.

To our knowledge, the relation of downward FACs and precipitating ions has not yet been investigated extensively. The electron energy flux is almost ten times larger than the ion (mostly proton) energy flux, and the contribution to the FACs from precipitating ions has usually been less discussed. Even though energetic protons are not an overall dominant energy source in the high-latitude region, they can be important at given location and time, for example near the equatorward boundary of the auroral oval in the afternoon and pre-midnight sectors (Hubert et al. 2001). Carter et al. (2016) also found that the downward R2 currents are allocated quite well with the auroral UV emission.

![Fig. 9](image_url)

Fig. 9 The locations of the peaks of upward (red) and downward (blue) FACs, as well as the precipitation electrons (dot within black circle) and ions (cross within black circle) at dawn and dusk sides, and also in the two hemispheres.
caused by precipitating protons at the dusk side. Our results also show smaller displacement (~1°) of the peaks between ion energy flux and downward R2 current at dusk side, while at dawn side, the precipitating ions are allocated more at the transition region between R1 and R2 currents and the peaks of the downward FACs and the ion energy flux is most displaced. Waters et al. (2001) commented that the post-midnight diffuse aurora is not necessarily produced solely by precipitating electrons, as diffuse auroral is brighter in the aurora far-ultraviolet (FUV) emission that corresponds to the distribution of precipitating protons. Murphy et al. (2013) also found that downward FACs typically occur in regions of diffuse aurora, but they did not separate whether the diffuse aurora was caused by precipitating electrons or protons. Our result suggests that downward R1 current is affected by both precipitating ions and electrons at the dawn side.

During comparison between Korth et al. (2014) and Carter et al. (2016), and our results, different observational limitations should also be considered when interpreting the similarities and differences. Both Korth et al. (2014) and Carter et al. (2016) used FACs estimates from AMPERE, while the former used N2 LBH intensity to derive electron energy flux, and the latter used wideband imaging camera (WIC) for representing the energy level of precipitation electrons. Both are indirect measurements of particle precipitation. Auroral emissions in the N2 LBH mainly reflect short (140–150 nm) and long (165–180 nm) wavelength ranges, and additional uncertainty exists in the empirical function between LBH intensity and electron energy flux (Robinson et al. 2018). WIC observations are continuously sensitive in the wavelength range from 140 to 190 nm. In addition, both studies compared datasets derived from different missions, AMPERE/GUVI and AMPERE/IMAGE, respectively. As there is no overlap between the AMPERE and IMAGE missions, Carter et al. (2016) optimized their comparison by sorting the FACs and WIC emissions into bins under similar IMF orientations. However, differences in the temporal and spatial coverage of the different missions exist. The simultaneously FAC and particle precipitations measured from the same platform of DMSP overcome at least differences in temporal and spatial coverage.

We also want to point out that although earlier studies showed a satisfactory agreement between the IMAGE-FUV observation and in situ particle measurements from satellites like FAST (Fast Auroral Snapshot Explorer) (Frey et al. 2001), DMSP (Gérard et al. 2001), and NOAA (National Oceanic and Atmospheric Administration) (Coumans et al. 2002), the DMSP SSJ5 instrument measures precipitating electrons/ions with energy range from 30 eV to 30 keV. Precipitating particles outside of this range may also contribute to the upward/downward currents, but cannot be captured by DMSP. Coumans et al. (2002) found that the contribution of protons with energy greater than 30 keV can exceed that of protons with energy less than 30 keV, more often observed between 1600–2100 MLT (see their Figs. 11 and 12). In addition, the DMSP SSJ instrument measures only the downward precipitating particles, but not the upward particles, like ion outflow in the cusp region, is also expected to cause FAC signatures (e.g., Strangeway et al. 2000). Hence, any conclusive interpretation of differences between FACs and precipitating particles, as shown in our study, should consider that it was derived from measurements of downward particle precipitation within a limited energy range.

Summary
In this study, we have used simultaneous observations of magnetic field and particle precipitations from the DMSP F16, F17, and F18 satellites between 2010 and 2014, to investigate the spatial relation between large-scale FACs and precipitation particles. The main findings can be summarized as:

1. FACs derived from DMSP magnetic data have been compared with Swarm satellite observations and evaluated against earlier studies. Their amplitudes are similar and the MLAT versus MLT distributions show known dependence on the IMF \(B_y\) and \(B_z\) components, which confirms the reliability of the DMSP magnetic data set.
2. Electron and ion energy fluxes intensify and extend to lower latitudes for enhanced southward \(B_z\), which is similar to the behavior of FACs.
3. Under weak northward and under southward $B_z$ conditions, R2 current peaks, at both dusk and dawn sides and in both hemispheres, are found to be close to peaks of the particle energy fluxes. The particle fluxes are slightly shifted poleward to the R2 currents, in average about 1° or less. For the same IMF conditions, the locations of the R1 currents and the respective particle flux peaks are even more displaced. The particle fluxes are shifted equatorward of the R1 currents, in average about 2° or more. The largest displacement is found between the downward R1 current and ion flux peak at the dawn side. Here, the ion flux peak is located about 3.5° MLAT equatorward of the R1 FAC peak at both hemispheres. Our results suggest that R1 current at dawn, often occurring together with diffuse auroras, is affected by both precipitating ions and electrons.

4. There exist systematic differences in the location of peaks of particle energy flux and large-scale FACs. On average, the FAC peaks enclose the particle energy flux peaks in an auroral band at both dusk and dawn sides and both hemispheres.

5. The particle precipitation, respectively, the electron or ion energy flux, both at dawn and dusk and in both hemispheres maximizes in average near the mean R2 current peaks. The particle precipitation maxima closer to the R1 current peaks are lower in magnitude. This is opposite to the known feature that R1 currents are on average stronger than R2 currents.

Geophysical Observatory, CHAMP: Challenging Minisatellite Payload; MLAT: Magnetic latitude; EMPÈRE: Active Magnetosphere and Planetary Electrodynamics Responses Experiment; TIMED: Thermosphere, Ionosphere, and Mesosphere Energetics and Dynamics; LBH: Lyman–Birge–Hopfield; GUVI: Global Ultraviolet Imager; UV: Ultraviolet; IMAGE: Imager for Magnetopause-to-Auroral Global Exploration; SSM: Special sensor magnetometer; SSJ: Special sensor J; VFM: Vector field magnetometer; ASM: Absolute scalar magnetometer; NEC: North, east, center; FLV: Far-ultraviolet; FAST: Fast Auroral Snapshot Explorer; NOAA: National Oceanic and Atmospheric Administration.

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Authors’ contributions
CX analyzed the FAC and particle precipitation data from DMSP, as well as the Swarm FAC measurements. PA processed and calibrated the magnetic measurements from DMSP. JR derived the FAC data set from DMSP magnetic data. CX and CS discussed aspects of data analysis and the content of the study. CX wrote and led the manuscript to which all authors contributed. All authors read and approved the final manuscript.

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Availability of data and materials
The solar wind data are available at https://omniweb.gsfc.nasa.gov. The Swarm FAC data can be downloaded at https://swarm-diss.eo.esa.int. The DMSP FACs data are available at ftp://magftp.gfz-potsdam.de/DMSP/FAC/ and the particle precipitation data are available at https://satdat.ngdc.noaa.gov/dmsp/data/f18/ssj/.

Competing interests
The authors declare that they have no competing interests.

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Appendix
See Fig. 10.
Fig. 10  a The data coverage of DMSP orbits in the northern hemisphere under different IMF conditions. For most of the bins, effective data point number is over 1000. b As similar as a but for the DMSP data coverage in the southern hemisphere.
Fig. 10 (continued)
