Usefulness of continuous probability distributions of rates for modelling radionuclide biokinetics in humans and animals

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Modelling the biokinetics of radionuclide excretion or retention is important in nuclear medicine and following accidental/malicious radioactivity releases. Sums of discrete exponential decay rates are often used, but we hypothesized that continuous probability distributions (CPD) of decay rates can describe the data more parsimoniously and robustly. We tested this hypothesis on diverse human and animal data sets involving various radionuclides (including plutonium, strontium, caesium) measured in the laboratory and in regions contaminated by the Fukushima and Chernobyl nuclear accidents. We used four models on each data set: mono-exponential (ME) with one discrete decay rate, bi-exponential (BE) with two rates, gamma-exponential (GE) with a Gamma distribution of stretched-exponential rates, and power-decay (PD) with a Gamma distribution of power-decay rates. Information-theoretic model selection suggested that radionuclide biokinetics, e.g. for plutonium in humans, are often better described by CPD models like GE and PD, than by discrete rates (ME and BE). Extrapolation of models fitted to data at short times to longer times was frequently more robust for CPD formalisms. We suggest that using a set of several CPD and discrete-rate models, and comparing them by information-theoretic methods, is a promising strategy to enhance the analysis of radionuclide excretion and retention kinetics.

Mathematical models of radionuclide excretion and retention kinetics from living organisms are important in a variety of contexts. For example, they are needed to estimate radiation doses and health risks from medical (e.g. nuclear medicine procedures), accidental (e.g. nuclear power plant accidents like Chernobyl or Fukushima), malicious (e.g. terrorist attacks using radioactive materials) or occupational (e.g. nuclear industry workers) exposures resulting in radionuclide dispersal and/or incorporation into the body. Such models are also needed in application to organisms other than humans, e.g. when nuclear power plant accidents such as Fukushima cause radioactive contamination of fish and game animals that are used for human consumption.

Radionuclides undergo well-understood physical decay. Importantly, however, many other chemical, biological and ecological processes also affect the kinetics of their removal from living organisms. Whereas physical decay has an exponential time dependence, these other processes can be much more complex and result in non-exponential time patterns. Detailed models have been developed to address this complexity, e.g. human radionuclide biokinetics models1 and models for the uptake and turnover of radionuclides in ecosystems2.

Complex models, however, have some important limitations. When the number of modelled processes that operate on different time scales and often have non-linear dependences is large, and the number of model parameters is correspondingly large, the model can become difficult to solve and parameter estimates can have very large uncertainties3,4. The latter phenomenon is sometimes called model “sloppiness”5.

Here we investigated the possibilities of using simple models, with small numbers of adjustable parameters, to describe radionuclide biokinetics data. For this purpose, we developed two new simple models based on the concept of a continuous probability distribution (CPD) of decay rates. The first, abbreviated as gamma-exponential (GE), combined the stretched exponential function6 with a Gamma distribution of rates. The second, abbreviated as power-decay (PD), combined a simplified version of the stretched hyperbola with a Gamma distribution of rates6. Using the Akaike information criterion with sample size correction (AICc) and multimodel inference

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measured under laboratory conditions: (a) Concentrations of humanized melanin-binding 111In-labeled IgG anti-
Ursus thibetanus), (b) Asian black bears Sus scrofa accident in Ukraine15. (II) Plasma data measured under laboratory conditions: (a) Concentrations of humanized melanin-binding 111In-labeled IgG anti-
Sus scrofa), (c) sika deer (Cervus nippon), and (d) ocellate spot skate (Okamejoi kenojei). (V), 137Cs radio-
activity concentrations in wild boars (Sus scrofa) in the area contaminated by the Chernobyl nuclear power plant
in the area contaminated by the Chernobyl nuclear power plant accident in Ukraine15.

The main goal of this study was to conduct a proof of principle investigation of whether or not continuous
probability distribution models like GE and PD could be reasonable for application in the fields of radionuclide
biokinetics and radioecology, compared with models with discrete rates like ME and BE. The concept of con-
tinuous rate distribution models would be potentially useful to investigators in these fields because different
types of biological and/or ecological processes can be summarized, producing robust numerical predictions,
by continuous distributions of rates. In situations when the kinetics data for a particular radionuclide are well
fitted by GE or PD functions, but not by ME or BE functions, the result suggests that many (rather than only 1 or
2) compartments and/or rates are likely to be involved. Detailed mechanistic models (e.g.) would be needed to
refine the information on specific rates and compartments, but such models may not be available for some studied
radionuclides and/or organisms.

Methods

Data sets. Our goal was to apply the proposed set of models (ME, BE, GE and PD) to a diverse set of real data
examples on radionuclide biokinetics from both humans and animals to assess a realistic range of model perform-
ance patterns, e.g. dominance of discrete rate models, dominance of CPD models, or no clearly best-supported
model(s). For this purpose, we identified the following data sets.

(I) Urinary excretion of plutonium in healthy human volunteers over time after administration12. Each
study subject ingested 234Pu citrate solution and was injected intravenously with additional 234Pu several
months later. This type of administration mimics a complicated pattern of plutonium incorporation into
the human body that can result from accidental or malicious radioactive contamination. Here we ana-
lysed the urinary excretion data on 5 male volunteers over time after injection, which were presented in
Table 4 of Ham et al.12. On this and all other data sets analysed here, we In-transformed the data to bring
the uncertainty distribution closer to Normal. We analysed the data from each subject separately, instead of
resorting to mixed-effects modelling, because we focused on comparing the performances of different
models by information-theoretic methods (described below), rather than on variations in parameter values
for a single model across subjects.

(II) Plasma concentrations of strontium in healthy human volunteers over time after administration13. Each
study subject ingested 89Sr and was simultaneously injected intravenously with 89Sr. We analysed the data
on 3 volunteers (separately for each one), which were presented in Table 4 of Li et al.13.

(III) Animal data measured under laboratory conditions. The first example involves a vertebrate model organ-
ism, the C57B6 female laboratory mouse. We used our own unpublished data on the kinetics of 111In-la-
beled IgG antibody to melanin in mouse blood over time after injection, corrected for physical decay of
111In. The data are presented in Supplementary Data File 1. The second example involves an invertebrate
model organism, the sea urchin Stronglylocentrotus nudus14. We analysed the data on 137Cs excretion in sea
urchins, which were obtained by first holding the urchins in radioactively contaminated water for 7 days,
and then transferring them to clean water and monitoring their radioactivity levels over time. These data
were presented in Fig. 1 of Nakamura et al.14, where each data point is the average value for 5 urchins. We
digitized the data using GetData Graph Digitizer 2.26 software (http://getdata-graph-digitizer.com/).

(IV) Radiocaesium (combined 137Cs and 134Cs) radioactivity concentrations in wild game animals inhabiting
the area contaminated by the Fukushima nuclear power plant accident in Japan: (a) wild boars (Sus scrofa),
(b) Asian black bears (Ursus thibetanus), and (c) sika deer (Cervus nippon) (Supplementary Data File 2).
These data were obtained from the Japan Atomic Energy Agency (JAEA) Radiation Monitoring Survey
Results of Wild Birds and Animals (https://emdb.jaea.go.jp/emdb/en/portals/1040501000/). We analysed the
data collected after September 1, 2011, i.e. approximately 6 months after the accident, because by this
time the accumulation of radionuclides within animals probably reached its full extent and their subse-
quent excretion kinetics could be evaluated. This time frame was chosen because earlier times would likely
be too soon for the ecological factors of radionuclide uptake and excretion in animals inhabiting the con-
taminated zone to reach a semi-equilibrium. Instances when the radiocaesium level was below detection
(usually <10 Bq/kg) were rare (0.8% of observations for wild boars, 4.1% for black bears, and 6.1% for sika
der), and we excluded them from analysis. The exclusion of these small fractions of very low radioactivity
measurements should not substantially bias the analysis results.

Since only the sum of 137Cs and 134Cs activities, rather than each isotope individually, was recorded in
animal tissues, we estimated the ratio of these two isotopes using data from the JAEA Unmanned Heli-
ocopter Monitoring in the Distribution Survey of Radioactive Substances (https://emdb.jaea.go.jp/emdb/
en/portals/b1020202/). We assumed that the isotope ratio would be the same for all studied animals, and
estimated it to be about 47.4% $^{134}\text{Cs}$ vs 52.6% $^{137}\text{Cs}$ on September 1, 2011. The ratio changed at subsequent times based on the physical half-lives of $^{134}\text{Cs}$ and $^{137}\text{Cs}$.

In addition to these data on terrestrial mammals, we also analysed data on $^{137}\text{Cs}$ in a bottom-dwelling marine fish, the ocellate spot skate ($\text{Okamejei kenojei}$) (Supplementary Data File 2). We used data collected in Fukushima prefecture, starting on April 1, 2012 (no earlier measurements were available), taken from the Results of the monitoring on radioactivity level in fisheries products database (http://www.jfa.maff.go.jp/e/inspection/). Because the fraction of instances with undetectable $^{137}\text{Cs}$ levels increased over time after the accident, the end date of May 13, 2015 was chosen to keep this fraction under 10%, so that exclusion of undetectable measurements from analysis would not substantially bias the results.

(V) $^{137}\text{Cs}$ radioactivity concentrations in wild boars ($\text{Sus scrofa}$) in the area contaminated by the Chernobyl nuclear power plant accident in Ukraine. These data were presented in Table 1 of Gulakov15. We used time since 1991, in years, as the independent variable for this analysis.

Because the focus of the current study is to investigate the time dependence of radioactivity excretion or retention processes, in all of these data sets we did not include other variables such as location. The effects of these other variables were not explicitly modelled here and were treated as components of random noise.

Models. The simplest mono-exponential (ME) decay model is represented by the following equation, where $R_{\text{ME}}(t)$ is the radioactivity at time $t$, $Q$ is the intercept parameter ($\exp[Q]$ is the radioactivity at $t = 0$), $PD(t)$ is the physical decay function of the radioactivity, and $k$ represents other radioactivity excretion and retention processes (e.g. biochemical, ecological):

$$R_{\text{ME}}(t) = \exp[Q] \times \exp[-k \times t] \times PD(t)$$

(1)

On a logarithmic scale, equation 1 can be rewritten as follows:

$$\ln[R_{\text{ME}}(t)] = Q + \ln[PD(t)] - k \times t$$

(2)

The physical decay function $PD(t)$ for $X$ radionuclides is the following sum of exponential dependences, where $C_j$ is the fractional contribution of the $j$-th radionuclide at $t = 0$, and $Th_j$ is the physical half-life of this radionuclide:

$$PD(t) = \sum_{j=1}^{X} C_j \times \exp\left[-\ln[2] \times \frac{t}{Th_j}\right]$$

(3)

The bi-exponential (BE) decay model is represented by the following equation, where $Q$ and $PD(t)$ have the same meanings as in the ME model, $k_2$ is the slow radioactivity reduction process, and $k_1$ is the additional reduction process, which acts on fraction $F$ of the radioactivity:

$$R_{\text{BE}}(t) = \exp[Q] \times (F \times \exp[-(k_1 + k_2) \times t] + (1 - F) \times \exp[-k_2 \times t]) \times PD(t)$$

(4)
The parametrization of equation 4 was chosen so that when \( k_1 \geq 0 \) and \( k_2 \geq 0 \), then \( k_1 + k_2 \geq k_3 \), thereby unambiguously defining \( F \) as the “fast-decaying” fraction of radioactivity and \( 1 - F \) as the “slow-decaying” fraction. On a logarithmic scale, equation 4 can be rewritten as follows:

\[
\ln[R_{\text{GE}}(t)] = Q + \ln[PD(t)] + \ln[F \times \exp[-(k_1 + k_2) \times t] + (1 - F) \times \exp[-k_3 \times t]]
\]

Of course, if either \( k_1 \) or \( k_2 \) approach 0, and/or if \( F \) approaches 0 or 1, the BE model simplifies to the ME model.

The proposed gamma-exponential (GE) decay model is based on the following assumptions: (1) The radioactive activity reduction processes, except physical decay, are summarized by the stretched exponential dependence \( \nu \). The same phenomenon occurs when \( \nu \) approaches 0, thereby unambiguously defining \( R \) as the “fast-decaying” fraction of radioactivity and \( 1 - R \) as the “slow-decaying” fraction. On a logarithmic scale, equation 4 can be rewritten as follows:

\[
R_{\text{GE}}(t) = \exp[Q \times \left( \int_{u=0}^{\infty} \exp[-u \times t'] \times P(u) \times du \right) \times PD(t)]
\]

The selected probability distribution \( P(u) \) is the following customized version of the Gamma distribution, where \( \mu \) is an adjustable parameter (the mean), \( \nu \) is a variance-determining parameter (so that variance \( = \nu \times \mu^2 \)), and \( \Gamma \) is the Gamma function:

\[
P(u) = \left[ u^{\nu-1} \times \exp \left[ -\frac{u}{\mu \times \nu} \right] / (\mu \times \nu)^{\nu/2} \right] \times \Gamma(1/\nu)
\]

Explicit solutions to equations 6 and 7 on the linear and logarithmic scales are as follows:

\[
R_{\text{GE}}(t) = \exp[Q \times PD(t) \times (1 + \mu \times \nu \times t')^{-(\nu/2)}],
\]

\[
\ln[R_{\text{GE}}(t)] = Q + \ln[PD(t)] - (1/\nu) \times \ln[1 + \mu \times \nu \times t']
\]

When parameter \( \nu \) approaches 0, the “width” of the \( P(u) \) distribution shrinks, and the GE model approaches the stretched exponential behaviour where \( \ln[R_{\text{GE}}(t)] \) depends on \( -\mu \times t' \). The same phenomenon occurs when \( \mu \) approaches 0. When parameter \( r \) approaches 1, the GE model approaches mono-exponential behaviour and becomes identical to the ME model (equation 1) where \( \mu \) is substituted for \( k_3 \).

The proposed power-decay (PD) model is based on the observation that deviations from simple exponential decay, such as the weighted sum of several exponential functions, can be parsimoniously represented by a power function like \( (1 + z \times t)^{-n} \), where \( z \) and \( n \) are adjustable parameters. The resulting PD model equation is as follows, where \( Q, P(u) \) and \( PD(t) \) have the same meanings as in equations above:

\[
R_{\text{PD}}(t) = \exp[Q \times \left( \int_{u=0}^{\infty} (1 + z \times t)^{-u} \times P(u) \times du \right) \times PD(t)]
\]
Explicit solutions to equation 9 on the linear and logarithmic scales are as follows:

\[
R_{pd}(t) = \exp[Q \times PD(t) \times (1 + \mu \times \nu \times \ln[1 + z \times t])^{-\frac{1}{2}}],
\]

\[
\ln[R_{pd}(t)] = Q + \ln[PD(t)] - (1/2) \times \ln[1 + \mu \times \nu \times \ln[1 + z \times t]]
\]

(10)

when parameter \( \nu \) approaches 0, \( \ln[R_{pd}(t)] \) depends on \(-\mu \times \ln[1 + z \times t] \). When parameter \( z \) approaches 0, the PD model approaches mono-exponential behaviour and becomes identical to the ME model (equation 1) where \( \mu \times z \) is substituted for \( k \).

Model fitting and parameter estimation. We used maximum likelihood estimation to find best-fit parameters for each of the 4 tested models (ME, BE, GE and PD) on each analysed data set. We maximized the following log-likelihood function^8, which assumes that the statistical uncertainties are normally distributed on a logarithmic scale, where \( N \) is the number of data points, \( \sigma \) is the statistical uncertainty magnitude parameter, \( W_i \) is the uncertainty "weight" for the \( i \)-th data point (so that the variance = \([\sigma \times W_i]^2 \)), \( R_{O,i} \) is the observed radioactivity at the \( i \)-th data point, and \( R_{M,i} \) is the corresponding radioactivity predicted by the M-th model (i.e. by one of the tested models):

\[
LL_{M} = -N \times \frac{[\ln(2 \times \pi) - \ln[\sigma]]}{2} - \sum_{i=1}^{N} \ln[W_i] - \frac{1}{2} \times \frac{1}{\sigma^2} \times \sum_{i=1}^{N} (\ln[R_{O,i}] - \ln[R_{M,i}]^2)^2/W_i^2
\]

(11)

For all analysed data sets, expect for the wild boar data from Chernobyl (data set V)\(^{14} \), the \( W_i \) values were unknown and presumably equal, so we set them to 1 for all data points. For the Chernobyl boar data, the radioactivity measurements in individual animals were not reported\(^{14} \), so we used the reported summary data: number of animals \( n_i \) mean \((AV_i)\) and maximum \((MAX_i)\) concentrations of \(^{137}\)Cs in boar muscle tissue (kBq/kg) at each \( i \)-th time point (during each studied year). These summary data were used to approximate \( \ln[R_{O,i}] \) and \( W_i \) as follows, where \( GE(n_i) \) is the median extreme value for a sample of size \( n_i \) taken from the standard Normal distribution:

\[
\ln[R_{O,i}] = \ln[AV_i], \quad W_i = (\ln[MAX_i] - \ln[AV_i])/GE(n_i)
\]

(12)

The values of \( GE(n_i) \) were estimated from 30,000 Monte Carlo simulations using Maple 2017\(^{®} \) software. For example, for \( n_i = 5 \) random values drawn from the standard Normal distribution, \( GE(n_i) = 1.13 \), which means that the median value (across all simulations) for the maximum of the sample is approximately 1.13 standard deviation units above the mean. Of course, \( GE(n_i) \) increases with \( n_i \) and for \( n_i = 20, GE(n_i) \) reached 1.82. Here we used \( GE(n_i) \) to approximate \( W_i \). Sensitivity calculations showed that the analysis results for this data set were not very sensitive to perturbations of \( W_i \).

The log likelihood function (equation 11) for each model on each data set was maximized using the sequential quadratic programming (SQP) algorithm in Maple 2017\(^{®} \) software\(^{17} \). The following procedure was used to find these 500 starting parameter combinations, and the combination which produced the highest log likelihood was retained as the best-fit parameter set. During the optimization procedure, the parameters were very sensitive to perturbations of \( \sigma \).

The number of adjustable parameters and \( LL_{M} \) is the maximized log-likelihood value:

\[
AICc_M = -2 \times LL_{M} + 2 \times \Lambda_M + 2 \times \Lambda_M \times (\Lambda_M + 1)/(N - \Lambda_M - 1)
\]

(13)

The model that achieves the lowest \( AICc \) value is considered to be best supported among those considered. The relative likelihood of the \( M \)-th model is called the evidence ratio (\( ER_M \)) and can be expressed as follows:

\[
ER_M = \exp \left\{-\frac{\Delta AICc_M}{2}\right\} \Delta AICc_M = AICc_M - AICc_{min}
\]

(14)

here, \( AICc_{min} \) is the lowest \( AICc \) value generated by the set of compared models.

The evidence ratio for \( M \)-th model, divided by the sum of the evidence ratios for all models, is the Akaike weight, \( W_M \). It represents the probability that the \( M \)-th model would be considered best-supported, among those tested, upon repeated sampling of the data. \( W_M \) is described by the following equation:

\[
W_M = ER_M / \sum_M ER_M
\]

(15)

As noted above, the BE, GE and PD models can all approach simple exponential behaviour (i.e. simplify to the ME model) if certain parameters approach limiting values such as zero. Consequently, if any of these models failed
to achieve a higher log likelihood than the ME model, their Akaike weights were assigned to the ME model. For example, if all 4 models had the same log likelihoods on a given data set, the simplest ME model was assigned an Akaike weight of 1 and all of the other more complex models were assigned weights of zero.

**Extrapolation of models from short to long times.** On each data set, we compared how the different models performed when each model was fitted to the data at short times and then extrapolated to data at longer times. The comparison was performed by calculating root mean squared error (RMSE) for each model on each data set under the following conditions: (1) Fitting the $M$-th model to all the data over the entire available time range, but calculating RMSE only on 1/3 of the data at the longest times. This produced $\text{RMSE}_{\text{Full,M}}$. (2) Fitting the $M$-th model only to those data points over the shortest 2/3 of the time range, but again calculating RMSE only on 1/3 of the data at the longest times. This produced $\text{RMSE}_{\text{Short,M}}$. Based on these two metrics, we calculated the RMSE loss ($\text{RMSE}_{\text{L,M}}$), which represents a difference in RMSE between conditions 1 and 2:

$$\text{RMSE}_{\text{L,M}} = \text{RMSE}_{\text{Short,M}} - \text{RMSE}_{\text{Full,M}}$$

Usually (but not always) $\text{RMSE}_{\text{L,M}}$ was a positive number, which suggests that RMSE calculated on data at long times was generally worse when the model was fitted to data at short times only instead of to all the data. The model that achieved the smallest RMSE loss, compared with other tested models, was considered to perform the best when extrapolated from short to long times. In other words, a small RMSE loss value indicates that the model has the right “shape” to decently predict the data at long times based on best-fit parameters obtained from fitting the data at short times. In addition, we also calculated RMSE for each model on each full data set.

**Results**

**Comparison of discrete rate model and CPD model behaviours.** Models that consist of sums of several discrete exponential rates are frequently used to analyse complex patterns of radioactivity excretion or reten-

### Analysis of data set 1.

Our analysis of real data on plutonium excretion in human volunteers (data set I) showed that the PD model strongly outperformed all other tested models (ME, BE and GE) on data from 4 out of 5 subjects (Table 1, Fig. 3, Supplementary Data File 3). This conclusion was supported by the PD model’s very high Akaike weight (essentially unity) and low root mean squared error (RMSE), compared with other models (Supplementary Data File 3). The PD model also outperformed other models when each model was fitted to data at short times and extrapolated to longer times. This is shown by the PD model’s low RMSE loss value (see Methods section), compared with other models (Supplementary Data File 3).

Subject 4 was an exception: the BE model was best-supported, whereas the PD model and the PD and GE models taken together had much lower Akaike weights (Table 1, Supplementary Data File 3). However, the best
RMSE loss was still achieved by the PD model (Table 1, Supplementary Data File 3). The likely reason for why the analysis results for subject 4 differed from those for the other subjects is the absence of data points at times between 7 and 84 days (Fig. 3).

Overall, these results suggest that the PD formalism, which assumes a continuous distribution of decay rates, reasonably approximated most of the plutonium data in a parsimonious manner, outperforming other tested models. The PD model (and the other formalisms tested here) are intended to generate simple and parsimonious “summaries” of the complex processes of plutonium biokinetics in the human body processes using a minimum number of parameters. In contrast, there is a different class of models - detailed and highly parametrized mechanistic formalisms, such as the model by Leggett et al.1. We believe that detailed biokinetics models and the simple
Summary models proposed here have different purposes. Both model classes can be applied to plutonium data, depending on the goals of the investigator.

Analysis of data set II. A different pattern was seen in the analysis of data set II on strontium retention in healthy human volunteers (Table 1, Fig. 4, Supplementary Data File 3). The decay kinetics appeared simpler than the kinetics for plutonium, possibly because of smaller data set size and shorter investigated time range. The ME model was best-supported on data for 2 out of the 3 subjects, but continuous rate distribution models achieved the best RMSE loss on data from all 3 subjects (Table 1, Supplementary Data File 3). These results suggest that strontium retention in the human body, perhaps expectedly, may deviate somewhat from ME kinetics, but that this deviation is difficult to quantify using the data and models described here.

Analysis of data set III. Data set III contains two examples of radioactivity retention kinetics in animals measured under laboratory conditions. The first example, melanin-binding 111In-labeled antibodies in mouse blood, was visually well described by the BE, GE and PD models (Fig. 5A), but because of the data set’s small size the simplest ME model achieved the highest Akaike weight (Supplementary Data File 3). The Akaike weights and RMSE values showed little difference between different models (Supplementary Data File 3 online), suggesting that none of these models clearly outperformed the others. RMSE loss was lowest for the PD model, which provides some indication that this formalism performed better than others when extrapolated from short to long times on this data set (Table 1, Supplementary Data File 3).

![Figure 4](https://example.com/figure4.jpg)  
Figure 4. Model best fits to the human strontium data (black circles), data set II. The y-axis represents ln-transformed plasma concentrations (µg/L) of 84Sr and 86Sr for three individuals, as a function of time after administration. The x-axis represents time after strontium administration. Green curve = ME model, cyan curve = BE model, blue curve = GE model, red curve = PD model.

![Figure 5](https://example.com/figure5.jpg)  
Figure 5. Model fits to data set III: laboratory animal data. (A) Model best fits to the data on melanin-binding 111In-labeled antibody in mouse blood (black circles). The y-axis represents percentage of injected radioactive dose per gram tissue, corrected for physical decay of 111In and ln-transformed. The x-axis represents time after injection. (B) Model best fits to radiocaesium data in sea urchins (black circles). The y-axis represents the 137Cs radioactivity concentration in the sea urchins, relative to its concentration in water, divided by the value at day 0 and ln-transformed. The x-axis represents time after the urchins were placed in non-radioactive water. In both panels, green curve = ME model, cyan curve = BE model, blue curve = GE model, red curve = PD model.
The second example, radiocaesium retention in sea urchins, appeared to be best described by the BE model (Table 1, Supplementary Data File 3), as reported in the original publication from which data were extracted. The GE and PD models, however, were not far behind the BE model, the sum of Akaike weights for the GE and PD models was >0.5 (Table 1), and all three formalisms produced similar best-fit curves (Fig. 5B). The BE, GE and PD models all had similar RMSE values, and the PD model had the best RMSE loss performance (Table 1, Supplementary Data File 3). The ME model had very poor support (Supplementary Data File 3). These results suggest that radiocaesium biokinetics in sea urchins involves at least two, and possibly more, decay rates. These rates are relatively rapid, which is consistent with observations that radiocaesium levels in sea urchins after the Fukushima nuclear disaster declined to undetectable levels within 3–4 years after the accident, whereas they remained elevated for much longer time in some other marine organisms like bottom-dwelling fish.

Analysis of data set IV. Data set IV contains four examples of ecological radioactivity data in terrestrial and aquatic animals in regions contaminated by the Fukushima nuclear power plant accident in Japan. On the first example, radiocaesium data in wild boars, the PD model simplified to approach the ME model, and the BE model exhibited unreasonable behaviour at short times (Table 1, Fig. 6A, Supplementary Data File 3). The GE model strongly outperformed all other tested formalisms (Table 1, Supplementary Data File 3). The ME model had very poor support (Supplementary Data File 3). These results suggest that radiocaesium biokinetics in sea urchins involves at least two, and possibly more, decay rates. These rates are relatively rapid, which is consistent with observations that radiocaesium levels in sea urchins after the Fukushima nuclear power plant accident declined to undetectable levels within 3–4 years after the accident, whereas they remained elevated for much longer time in some other marine organisms like bottom-dwelling fish.

In contrast, radiocaesium in black bears from the same contaminated area of Japan appeared to decrease over time with simpler kinetics than in wild boars: the ME model achieved the highest support on black bear data (Table 1, Supplementary Data File 3). The best fits from all other models essentially overlapped the one from the ME model (Fig. 6B), and the ME model was favoured in information theoretic analysis due to its simplicity. Sika deer data from the Fukushima-contaminated area displayed yet another pattern: they were best described by the BE model, which predicted rapid decay at short times and essentially no decay (other than physical) at longer times (Table 1, Fig. 6C, Supplementary Data File 3).

Model-averaged (MMI) best-fit curves for boars, bears and deer are shown in Fig. 6E. These species living in a contaminated habitat were subjected to continuous intakes of radionuclides and hence approached a semi-equilibrium in terms of radionuclide intake and excretion, which is often attained after a year or two after...
the initial contamination event. The time-patterns seen here therefore primarily reflect ecological processes, rather than the metabolic patterns in the species.

The last example, aquatic ocellate spot skate data, favoured the BE model, although the sum of Akaike weights for the GE and PD models was >0.5 (Table 1, Fig. 6D, Supplementary Data File 3). Most of the tested models produced visually similar fits to the data, and their RMSE and RMSE loss values were similar (Table 1, Fig. 6D, Supplementary Data File 3). The ocellate spot skate is an example of a bottom-dwelling fish species in which radiocaesium levels remained elevated even several years after the Fukushima accident19.

Analysis of data set V. Our analysis of the final data set (V) on 137Cs in wild boars from the area contaminated by the Chernobyl nuclear power plant accident15 suggested that caesium levels in Chernobyl boars (Table 1, Fig. 7) did not decrease at all due to ecological processes over the studied period. The best-fit curves from all tested models therefore overlapped, with the decrease in caesium radioactivity attributable exclusively to physical decay (Fig. 7). This pattern likely reflects the ecological turnover of radiocaesium and the dietary habits of the studied species. Wild boars consume mushrooms, which strongly accumulate caesium20, and the boar radioactivity fluctuations over the years may arise from the fact that particular years exhibit a more intense mushroom growth.

Discussion
We applied two proposed models with continuous probability distributions of decay rates (GE and PD) and two conventional models with sums of discrete decay rates (ME and BE) to a deliberately diverse collection of human and animal data sets on the biokinetics of radioactive materials in living organisms. This collection included kinetics data on multiple elements (plutonium, strontium, caesium, indium) generated in the laboratory (data sets I-III) and under natural conditions in areas contaminated by nuclear accidents (data sets IV-V). The focus of the study was to evaluate how the selected assemblage of four models behaves on each of these different example data sets, and what inferences these models can provide about the biological, chemical and ecological processes involved in radionuclide kinetics.

The rationale for developing and using these models was that simple formalisms can often bring out the key features of the studied system and produce reliable interpolations and extrapolations, based even on limited amounts of data, using only a few adjustable parameters3,4,21. In the case of radionuclide biokinetics, simple easy to use models are needed for predicting radioactivity values at long times based on data measured at shorter times. Such predictions are important in many situations, e.g. for estimating the cumulative radiation dose absorbed by a person exposed to radionuclides in a medical setting, or for estimating how long after a nuclear accident will game animals in the area retain radioactivity concentrations too high for human consumption.

The most common approach to simple modelling of radionuclide biokinetics involves a sum of exponential decay rates22–28. A certain fraction of radioactivity (e.g. the one contained in a given organ or body “compartment”) is assumed to be eliminated at one rate, another fraction at a different (e.g. slower) rate, and so on. In this system, the different rate constants have compensatory effects: a decrease in one rate can be compensated for by increases in other rate(s). Consequently, parameter uncertainties tend to be large4 and extrapolation of model predictions to times longer than those used for model fitting tends to be unreliable.

These problems can be addressed by assuming that the complex chemical, biological and ecological processes during radionuclide excretion or retention can be approximated by a continuous distribution of decay rates. In other words, instead of using a sum of several discrete rates, one can use an integral over an infinite number of
rates described by a certain probability distribution. This approach, which is known in pharmacokinetics, chemistry and physics, allows multiple incompletely understood processes to be combined into a “mixture” represented by only a few adjustable parameters. Specifically, solutions for kinetic equations that characterize relaxation in disordered systems often lead to simple functions such as the stretched exponential or stretched hyperbola, which are known since the 19th century.

We used this line of reasoning to develop the GE and PD models. The diverse data sets chosen for model testing were intended to represent a variety of scenarios where modelling of radionuclide biokinetics can be important. They also represent a variety of decay patterns, from simple to more complex. We fitted each model to each data set and compared model performances by information theoretic approaches. We also compared the ability of each model on each data set to extrapolate from short times to longer times.

As expected, the relative performances of different models varied widely depending on data set. The variability in model performance on different data sets likely reflects underlying differences in radionuclide biokinetics due to differences in organism, radionuclide type and environment, and the effects of data set size and structure. Whereas on some data sets one model was dominant (e.g. the PD model on data set I), in other cases no model was strongly favoured over others (e.g. on data set III). On one data set (V) none of the models was required to explain the data because the time trend was sufficiently explained by physical radionuclide decay alone. Larger data sets (e.g. IV, I) of course tend to support more complex models and provide more information for model selection, than smaller ones (e.g. III, VII).

The sum of Akaike weights for the proposed GE and PD models can be interpreted as evidence for >2 decay rates being involved. This sum was frequently >0.5 among the human and animal data sets analysed here (Table 1). Even more frequently, the GE and/or PD models achieved the lowest RMSE loss, suggesting that these models tend to be more robust than discrete rate models (ME and BE) during extrapolation from short to long times (Table 1). Even greater robustness can potentially be achieved by model averaging (MMI), particularly when no single model has a clearly dominant Akaike weight.

Taken together, our results suggest that radionuclide biokinetics in living organisms are often sufficiently complex to be reasonably described by a continuous distribution of decay rates (as assumed in the GE and PD models), than by the commonly used assumption of one or two discrete rates (as in the ME and BE models). For example, the PD model showed strong performance on human plutonium excretion data and is notably much more parsimonious than the extensively parametrized multi-compartment kinetic models. Consequently, we think that CPD models are useful alternatives to discrete rate models.

The CPD approach can be used to determine whether the radionuclide biokinetics involve multiple compartments and/or rate constants. Robust results based on all tested models can be obtained using MMI, and they can prove to be very useful in many situations involving medical use of radionuclides and accidental or malicious radioactive contamination.

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Author Contributions
I.S. and E.D. conceived the study and prepared the manuscript. I.S. carried out the analysis. All authors discussed the model implementation and results, and contributed to editing the manuscript.

Additional Information
Supplementary information accompanies this paper at https://doi.org/10.1038/s41598-018-38046-9.

Competing Interests: The authors declare no competing interests.

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