Photon spectra from quark generation by WIMPs

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Abstract. If the present dark matter (DM) in the Universe annihilates into Standard Model (SM) particles, it must contribute to the gamma ray fluxes that are detected on the Earth. The magnitude of such contribution depends on the particular DM candidate, but certain features of these spectra may be analyzed in a model-independent fashion. In this work we provide the fitting formula valid for the simulated photon spectra from WIMP annihilation into light quark-anti quark ($q\bar{q}$) channels in a wide range of WIMP masses. We illustrate our results for the $c\bar{c}$ channel.

Keywords: Dark matter, indirect searches, WIMPs, photon spectra and quark pairs annihilation.

PACS: 95.35.+d 98.80.Cq.

I. INTRODUCTION

According to present observations of large scale structures, CMB anisotropies and light nuclei abundances, DM cannot be accommodated within the SM of elementary particles. Indeed, DM presence is a required component on cosmological scales, but also to provide a satisfactory description of rotational speeds of galaxies, orbital velocities of galaxies in clusters, gravitational lensing of background objects by galaxy clusters and the temperature distribution of hot gas in galaxies and clusters of galaxies. The experimental determination of the DM nature will require the interplay of collider experiments and astrophysical observations. These searches use to be classified in direct or indirect searches (see [1] and references in Introduction in [2]). Concerning direct ones, the elastic scattering of DM particles from nuclei should lead directly to observable nuclear recoil signatures although the weak interactions between DM and the standard matter makes DM direct detection extremely difficult.

On the other hand, DM might be detected indirectly, by observing their annihilation products into standard model particles. Thus, even if WIMPs (Weakly Interacting Massive Particles) are stable, two of them may annihilate into ordinary matter such as quarks, leptons and gauge bosons. Their annihilation in different places (galactic halo, Sun, etc.) produce cosmic rays to be discriminated through distinctive signatures from the background. After WIMPs annihilation a cascade process occurs. In the end the stable particles: neutrinos, gamma rays, antimatter... may be observed through different devices. Neutrinos and gamma rays have the advantage of maintaining their original direction due to their null electric charges.

This communication precisely focuses on photon production coming from $q\bar{q}$ channels (except $t\bar{t}$ channel). Photon fluxes in specific DM models are usually obtained by software packages such as DarkSUSY and micrOMEGAs based on PYTHIA Monte Carlo event generator [3] after having fixed a WIMP mass for the particular SUSY model under consideration. In this sense, the aim of this investigation is to provide fitting functions for the photon spectra corresponding to each individual annihilation $q\bar{q}$ channel and, in addition, determine the dependence of such spectra on the WIMP mass in a model independent way. This would allow to apply the results to alternative DM candidates for which software packages have not been developed. On the other hand, the information about channel contribution and mass dependence can be very useful in order to identify gamma-ray signals for specific WIMP candidates and may also provide relevant information about the photon energy distribution when $q\bar{q}$ pairs annihilate.

Let us remind that the $\gamma$-ray flux from the annihilation of two WIMPs of mass $M$ into two SM particles coming from all possible annihilation channels (labelled by the subindex $i$) is given by:

\[
\frac{d\Phi_{\gamma}^{DM}}{dE_{\gamma}} = \frac{1}{4\pi M^2} \sum_i \langle \sigma_i v_i \rangle \frac{dN_i^\gamma}{dE_{\gamma}} \times \frac{1}{\Delta\Omega} \int_{\Delta\Omega} d\Omega \int_{l.o.s.} \rho^2(r) ds , \tag{1}
\]

where $\langle \sigma_i v_i \rangle$ holds for the thermal averaged annihilation cross-section of two WIMPs into two $i$th channel SM particles and $\rho$ is the DM density. The integral is performed along the line of sight (l.o.s.) to the target and averaged over the detector solid angle $\Delta\Omega$. 

The IX International Conference on Quark Confinement and the Hadron Spectrum - QCHS IX
AIP Conf. Proc. 1343, 595-597 (2011); doi: 10.1063/1.3575105
© 2011 American Institute of Physics 978-0-7354-0899-9/$30.00
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II. PROCEDURE

We have used the particle physics PYTHIA software [3] to obtain our results. The WIMP annihilation is usually split into two separated processes: The first describes the annihilation of WIMPs and its SM output. The second one considers the evolution of the obtained SM unstable products. Due to the expected velocity dispersion of DM, we expect most of the annihilations to happen quasi-statically. This fact allows to state that by considering different center of mass (CM) energies for the obtained SM particles pairs from WIMP annihilation process, we are indeed studying different WIMP masses, i.e., $E_{\text{CM}} \simeq 2M$. The procedure to obtain the photon spectra is thus straightforward: For a given pair of SM particles which are produced in the WIMP annihilation, we count the number of photons in bins for the variable $x \equiv E_{\gamma}/M$.

Once the PYTHIA simulations have been performed, the required parametrization to fit the data for the $q\bar{q}$ channels (except $t\bar{t}$) may be written as:

$$\frac{dN_{\gamma}}{dx} = \frac{a_{1}}{x^{3}} \exp \left( -b_{1}x^{a_{1}} - b_{2}x^{a_{2}} - \frac{c_{1}}{x^{d_{1}}} + \frac{c_{2}}{x^{d_{2}}} \right) + q \ln[p(1-x)] \frac{x^{2} - 2x + 2}{x}$$

The parameters in expression (2) were considered to be WIMP mass dependent. After a fitting process they were determined for different WIMP masses, in a range varying from 50 to 7000 (or 8000) GeV. Mass dependences for the parameters in (2) were fitted by using power laws.

III. $c\bar{c}$ CHANNEL

In order to illustrate the explained procedure, we study the $c\bar{c}$ channel. For this channel there are five mass dependent parameters in expression (2): $b_{1}, n_{1}, c_{1}, d_{1}$ and $p$. Presented in Table I. The mass independent parameters are $a_{1}, b_{2}, n_{2}, c_{2}$ ($d_{2}$ is thus irrelevant) and $q$.

| $M$ (GeV) | $b_{1}$ | $n_{1}$ | $c_{1}$ | $d_{1}$ | $p$ |
|----------|--------|--------|--------|--------|-----|
| 50       | 5.93   | 2.35   | 0.239  | 0.428  | 210 |
| 100      | 5.48   | 2.08   | 0.283  | 0.374  | 379 |
| 200      | 4.98   | 1.86   | 0.330  | 0.330  | 673 |
| 500      | 4.50   | 1.65   | 0.378  | 0.288  | 1230|
| 1000     | 4.00   | 1.50   | 0.406  | 0.264  | 2110|
| 2000     | 3.70   | 1.35   | 0.432  | 0.245  | 4050|
| 5000     | 3.27   | 1.17   | 0.470  | 0.221  | 8080|
| 8000     | 3.08   | 1.11   | 0.494  | 0.208  | 12000|

II. CONCLUSIONS

In this work, we have studied the photon spectra coming from WIMP pair annihilation into $q\bar{q}$ pairs for all the channels (except $t\bar{t}$). The covered WIMP mass range was from 50 GeV to 8000 GeV. Simulated spectra covered the whole accessible energy interval: from extremely low energetic photons till photons with one half of the available total center of mass energy.

Once the spectra were simulated, an analytical expression (2) was proposed to fit the data. This expression depends on both WIMP mass dependent and independent parameters. Our results can both provide a better understanding of the DM annihilation channels into photons and save an important amount of unnecessary Monte Carlo simulations. This fact is particularly important for high energy photons, whose production rate is very suppressed.

Calculations for all $q\bar{q}$ channels are available at the website http://teorica.fis.ucm.es/~PaginaWeb/downloads.html. This work was partially supported by MULTIDARK CSD2009-00064.

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FIGURE 1. Photon spectra for four different WIMP masses (100, 200, 1000 and 5000 GeV) in the $c\bar{c}$ annihilation channel. Red dotted points are PYTHIA simulations and solid lines correspond to the proposed fitting functions.

FIGURE 2. Mass dependence of $c_1$ and $b_1$ parameters for $c\bar{c}$ annihilation channel. Crossed points are parameters values found after the fitting process for each WIMP mass and solid lines correspond to the proposed fitting functions.
