Comparison of fuzzy Adomian decomposition method with fuzzy VIM for solving fuzzy heat-like and wave-like equations with variable coefficients

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Abstract
In this work, we investigate the fuzzy Adomian decomposition method (ADM) and fuzzy variational iteration method (VIM) applied to solving fuzzy heat-like and wave-like equations with variable coefficients, in the sense of gH-differentiability. The methods clearly are very efficient and powerful techniques in finding the solutions to the proposed equations. We also illustrate them by some examples. The results reveal that the methods are very effective, convenient, and accurate mathematical tools for solving fuzzy heat-like and wave-like equations with variable coefficients.

Keywords: Fuzzy numbers; Fuzzy-valued functions; gH-differentiability; Fuzzy Adomian decomposition method; Fuzzy variational iteration method; Fuzzy heat-like and wave-like equations

1 Introduction
The fuzzy set theory is a powerful tool for modeling uncertainty and processing vague or subjective information in mathematical models, which have been applied to a wide variety of real problems; for example, see [4, 17, 19, 20, 56]. Especially, the fuzzy differential equations (FDEs) are a very important topic from the theoretical point of view; some authors are interest in these works, for example, Fei [21] studied the existence and uniqueness of solutions for fuzzy random differential equations (FRDEs) with non-Lipschitz coefficients. The concept of differential equations in the fuzzy environment was introduced by Kaleva [42] in 1986. Kaleva [43] also presented the Cauchy problem for FDEs. Lopez [44] obtained comparison results for FDEs. Friedman and Kandel [45] introduced numerical solutions of FDEs. Mizukoshi et al. [46] studied the FDEs and the extension principle, as well as the applied point of view [2, 3, 9]; see, for example, [14, 15]. The concept of the fuzzy derivative was initially proposed by Chang and Zadeh [16] in 1972. Dubois and Prade [18] introduced the extension principle. Other methods have been considered by Puri and Ralesu [51] and Goetschel and Voxman [23].

The fuzzy partial differential equations (FPDEs) attracted a great deal of attention among scientists and engineers because of their frequent involvement in the modeling of...
numerous industrialized applications, for example, electromagnetic fields, heat and mass transfer, meteorology, static and dynamic of structures, biomechanics, and other areas of science and engineering. The FPDEs have considered by many researchers. Allahviranloo [7] used the Adomian decomposition method (ADM) for solving the fuzzy heat equation. Alikhani and Bahrami [6] presented the FPDEs under the cross product of fuzzy numbers. Jameel et al. [41] proposed a new approximation method for solving fuzzy heat equations. Stepnicka and Valasek [54] introduced the numerical solution of partial differential equations (PDEs) with the help of a fuzzy transform. Jameel [39] studied the semianalytical solution of the heat equation in a fuzzy environment. Allahviranloo et al. [10] have presented the fuzzy solutions for heat equation based on generalized Hukuhara differentiability. Osman et al. [47] used the differential transform method for solving fuzzy fractional wave equations under generalized Taylor’s formula, gH-differentiability, and fuzzy Coputo’s derivative.

The identification of the Lagrange multiplier is extremely very important in mathematical tools and can be optimized by using the variational theory. He [33] introduced a generalized variational principle and obtained a parameterized generalized variational principle for large deformation analysis of circular cylinders by the semiinverse method. He and Ain [36] established new promises and future challenges of fractal calculus. from two-scale thermodynamics to fractal variational principle. He [34] used the variational principle and periodic solution of the Kundu–Mukherjee–Naskar equation.

The Adomian decomposition method (ADM) has been lately intensively considered by scientists and engineers and studied for solving linear, nonlinear differential, and integral problems. Hamoud and Ghadle [28] successfully applied the ADM and the modified Laplace ADM to find approximate solutions of a nonlinear fractional Volterra–Fredholm integro-differential equation under Caputo’s fractional derivative. Hamoud et al. [30] studied the behavior of approximate solutions of the Caputo fractional Volterra–Fredholm integro-differential equations. At the same time, several researchers used this method for solving many equations in the fuzzy environment; see [1, 13, 27, 48, 49, 53].

The variational iteration method (VIM) was originally proposed by He [32]. Later, several authors found VIM to be an effective way for solving some equations; see [5, 26, 29, 31, 35, 37, 38]. Furthermore, in a fuzzy environment, several authors studied the VIM. Allahviranloo et al. [8] considered the fuzzy VIM for solving nonlinear fuzzy differential equations. Jameel [40] presented an approximate analytical solution for an initial value problem involving the fuzzy Duffing ordinary differential equation by using VIM. Hamoud et al. [25] studied the ADM, VIM, and homotopy analysis method for solving fuzzy Volterra–Fredholm integral equations. Zhang et al. [57] proposed a fuzzy model for dye removal and suggested to study a transport model of the direct textile industry wastewater. In this paper, we present a comparison of the fuzzy Adomian decomposition method and the fuzzy variational iteration method for solving fuzzy heat-like and wave-like equations with variable coefficients under the gH-differentiability.

This work is structured as follows. In Sect. 2, we recall some definitions of fuzzy numbers, fuzzy-valued functions, and strongly gH-differentiable functions. In Sect. 3, we propose the analysis of fuzzy heat-like and wave-like equations with variable coefficients under gH-differentiability by using the fuzzy ADM and VIM. In Sect. 4, we provide some examples to show the efficiency and simplicity of the ways they were developed and derived. Finally, a conclusion is given in Sect. 5.
2 Basic concepts

The basic definition of fuzzy numbers is given in [58]. We denote the set of all real numbers by \( R \) and the set of all fuzzy numbers on \( R \) by \( E^1 \). A fuzzy number is a mapping \( u: R \to [0,1] \) with the following properties:

1. \( \tilde{u} \) is normal, that is, there exists \( x_0 \in R \) with \( \tilde{u}(x_0) = 1 \);
2. \( \tilde{u} \) is a convex fuzzy set (i.e., \( \tilde{u}(\lambda x + (1 - \lambda)y) \geq \min\{\tilde{u}(x), \tilde{u}(y)\}, \forall \lambda \in [0,1], x, y \in R \);
3. \( \tilde{u} \) is semicontinuous on \( R \);
4. \( \{x \in R : \tilde{u}(x) > 0\} \) is compact, where \( \overline{A} \) denotes the closure of \( A \).

**Definition 2.1** ([45, 52, 58]) A fuzzy number \( \tilde{u} \) in parametric form is a pair \((\underline{u}, \overline{u})\) of functions \( \underline{u}(r), \overline{u}(r), 0 \leq r \leq 1 \), satisfying the following requirements:

1. \( \underline{u} \) is a bounded increasing left-continuous function,
2. \( \overline{u} \) is a bounded decreasing left-continuous function,
3. \( \underline{u} \leq \overline{u}(r), 0 \leq r \leq 1 \).

Applying Zadeh's extension principle, the operation of addition on \( E^1 \) is defined by

\[
(\tilde{u} + \tilde{v})(x) = \sup_{y \in R} \min\{\mu(y), \nu(x-y)\}, \quad x \in R, \tag{2.1}
\]

and the scalar multiplication of fuzzy numbers is given by

\[
(k \odot \tilde{u})(x) = \begin{cases} 
\tilde{u}(x/k), & k > 0, \\
\tilde{0}, & k = 0,
\end{cases}
\]

where \( \tilde{0} \in E^1 \).

The Hausdorff distance between fuzzy numbers is the map \( d: E^1 \times E^1 \to [0, +\infty] \) given by

\[
d(\tilde{u}, \tilde{v}) = \sup_{r \in [0,1]} \max\{|\mu_r - \nu_r|, |\pi_r - \pi_r|\} \tag{2.2}
\]

for \( \tilde{u} = (\underline{u}, \overline{u}) \) and \( \tilde{v} = (\underline{v}, \overline{v}) \subset R \) [11]. It is easy to see that \( d \) is a metric in \( E^1 \) and has the following properties [50]:

1. \((E^1, D)\) is a complete metric space,
2. \( D(\tilde{u} \oplus \tilde{w}, \tilde{v} \oplus \tilde{w}) = D(\tilde{u}, \tilde{v}), \forall \tilde{u}, \tilde{v}, \tilde{w} \in E^1 \),
3. \( D(\tilde{u} \oplus \tilde{v}, \tilde{w} \oplus \tilde{w}) \leq D(\tilde{u}, \tilde{w}) + D(\tilde{v}, \tilde{w}), \forall \tilde{u}, \tilde{v}, \tilde{w}, \tilde{w} \in E^1 \),
4. \( D(k \odot \tilde{u}, k \odot \tilde{v}) = |k|D(\tilde{u}, \tilde{v}), \forall \tilde{u}, \tilde{v} \in E^1, k \in R \).

Note that functions \( \tilde{f}: A \to E^1, A \subseteq R \), are further called fuzzy-valued functions, whereas functions \( \tilde{f}: A \to R, A \subseteq R \), are called real-valued functions (or crisp functions). The \( r \)-level representation of fuzzy-valued function \( \tilde{f} \) can be expressed by \( \tilde{f}(x; r) = \{[\tilde{f}(x; r), \tilde{f}(x; r)] : x \in A \subseteq R \) and \( 0 \leq r \leq 1 \).

**Definition 2.2** ([42]) A fuzzy-valued function \( \tilde{f}: [a, b] \to E^1 \) is said to satisfy condition (gH) on \([a, b]\) if for any \( x_1 < x_2 \in [a, b] \), there exists \( \tilde{u} \in E^1 \) such that \( \tilde{f}(x_2) = \tilde{f}(x_1) + \tilde{u} \). We call \( \tilde{u} \) the gH-difference of \( \tilde{f}(x_2) \) and \( \tilde{f}(x_1) \), denoted \( \tilde{f}(x_2) \ominus_{gH} \tilde{f}(x_1) \). For brevity, we always assume that it satisfies condition (gH) when we deal with the subtraction of fuzzy numbers.
**Definition 2.3** ([12, 55]) The generalized Hukuhara difference of two fuzzy numbers \( \tilde{u}, \tilde{v} \in E^1 \) is defined as follows:

\[
\tilde{u} \otimes_{gH} \tilde{v} = \tilde{w} \iff \begin{cases} 
(i) & \tilde{u} = \tilde{v} \oplus \tilde{w}, \text{ or} \\
(ii) & \tilde{v} = \tilde{u} \oplus (-\tilde{w}). 
\end{cases} \tag{2.3}
\]

In terms of the \( r \)-levels, we get \( [\tilde{u} \otimes_{gH} \tilde{v}] = [\min\{\tilde{u}_r - \tilde{v}_r, \tilde{u}_r - \tilde{v}_r\}, \max\{\tilde{u}_r - \tilde{v}_r, \tilde{u}_r - \tilde{v}_r\}] \), and if the H-difference exists, then \( \tilde{u} \otimes \tilde{v} = \tilde{u} \otimes_{gH} \tilde{v} \), conditions for the existence of \( \tilde{w} = \tilde{u} \otimes_{gH} \tilde{v} \in E^1 \) are

Case (i)

\[
\begin{align*}
\tilde{w}_r &= u_r - v_r \quad \text{and} \\
\tilde{w}_r &= \bar{r}_r - \bar{v}_r, \\
\forall r &\in [0, 1], \\
\text{with} \tilde{w}_r &\text{ increasing,} \tilde{w}_r &\text{ decreasing,} \tilde{w}_r &\leq \bar{w}_r, \\
\end{align*} \tag{2.4}
\]

Case (ii)

\[
\begin{align*}
\tilde{w}_r &= \bar{u}_r - \bar{v}_r \quad \text{and} \\
\tilde{w}_r &= u_r - v_r, \\
\forall r &\in [0, 1], \\
\text{with} \tilde{w}_r &\text{ increasing,} \tilde{w}_r &\text{ decreasing,} \tilde{w}_r &\leq \bar{w}_r. \\
\end{align*} \tag{2.5}
\]

It is easy to show that (i) and (ii) are both valid if and only if \( \tilde{w} \) is a crisp number. In this case, it is possible that the gH-difference of two fuzzy numbers does not exist. To address this shortcoming, a new difference of fuzzy numbers was introduced in [12].

**Definition 2.4** ([11]) Let \( \tilde{u}(x, t) : D \to E^1 \) and \( (x_0, t) \in D \). We say that \( \tilde{u} \) is strongly generalized Hukuhara differentiable on \( (x_0, t) \) (gH-differentiable for short) if there exists an element \( \frac{\partial \tilde{u}}{\partial x}_{gH}(x_0, t) \in E^1 \) such that

(i) for all \( h > 0 \) sufficiently small, \( \tilde{u}(x_0 + h, t) \otimes_{gH} \tilde{u}(x_0, t) \) and the limits (in the metric \( D )

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0 + h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0, t) \otimes_{gH} \tilde{u}(x_0 - h, t)}{h} = \frac{\partial \tilde{u}}{\partial x}_{gH}(x_0, t),
\]

or

(ii) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0, t) \otimes_{gH} \tilde{u}(x_0, t) \) and the limits

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0 + h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0 - h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \frac{\partial \tilde{u}}{\partial x}_{gH}(x_0, t),
\]

or

(iii) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0 + h, t) \otimes_{gH} \tilde{u}(x_0, t) \) and the limits

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0 + h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0 - h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \frac{\partial \tilde{u}}{\partial x}_{gH}(x_0, t),
\]

or

(iv) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0, t) \otimes_{gH} \tilde{u}(x_0, t) \) and the limits

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0 + h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0 - h, t) \otimes_{gH} \tilde{u}(x_0, t)}{h} = \frac{\partial \tilde{u}}{\partial x}_{gH}(x_0, t).
\]
Definition 2.5 ([24]) Let \( \tilde{u}(x, t) : D \to E^1 \) and \((x_0, t) \in D \). We define the \( n \)th-order derivative of \( \tilde{u} \) as follows: we say that \( \tilde{u} \) is strongly generalized Hukuhara differentiable of the \( n \)th order at \((x_0, t)\) (\( g_H \)-differentiable for short) if there exists elements \( \frac{\partial^s \tilde{u}}{\partial x^s} |_{(x_0, t)} \in E^1 \), \( s = 1, 2, \ldots, n \), such that

(i) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0 + h, t) \odot g_H \tilde{u}(x_0, t) \), \( \tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 - h, t) \) and the limits (in the metric \( D \))

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0 + h, t) \odot g_H \tilde{u}(x_0, t)}{h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0 - h, t) \odot g_H \tilde{u}(x_0, t)}{h} = \frac{\partial^s \tilde{u}}{\partial x^s} |_{(x_0, t)},
\]

or

(ii) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 + h, t) \), \( \tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 - h, t) \) and the limits

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 + h, t)}{-h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 - h, t)}{-h} = \frac{\partial^s \tilde{u}}{\partial x^s} |_{(x_0, t)},
\]

or

(iii) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0 + h, t) \odot g_H \tilde{u}(x_0, t) \), \( \tilde{u}(x_0 - h, t) \odot g_H \tilde{u}(x_0, t) \) and the limits

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0 + h, t) \odot g_H \tilde{u}(x_0, t)}{h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0 - h, t) \odot g_H \tilde{u}(x_0, t)}{h} = \frac{\partial^s \tilde{u}}{\partial x^s} |_{(x_0, t)},
\]

or

(iv) for all \( h > 0 \) sufficiently small, there exist \( \tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 + h, t) \), \( \tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 - h, t) \) and the limits

\[
\lim_{h \to 0^+} \frac{\tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 + h, t)}{-h} = \lim_{h \to 0^+} \frac{\tilde{u}(x_0, t) \odot g_H \tilde{u}(x_0 - h, t)}{h} = \frac{\partial^s \tilde{u}}{\partial x^s} |_{(x_0, t)}.
\]

3 Fuzzy heat-like and wave-like equations with variable coefficients

In this section, we analyze fuzzy heat-like and wave-like equations with variable coefficients under \( g_H \)-differentiability by using the fuzzy Adomian decomposition method and fuzzy variational iteration method as follows.

First, we consider the fuzzy heat-like equation of the form

\[
\tilde{u}_t = f(x, y, z) \odot \tilde{u}_{xx} \odot g(x, y, z) \odot \tilde{u}_{yy} \odot h(x, y, z) \odot \tilde{u}_{zz},
\]

\[
0 < x < a, 0 < y < b, 0 < z < c, t > 0, f(x, y, z) \geq 0, g(x, y, z) \geq 0, h(x, y, z) \geq 0,
\]
with the Neumann boundary conditions
\[
\tilde{u}_s(0, y, z, t) = \tilde{f}_1(y, z, t), \quad \tilde{u}_s(a, y, z, t) = \tilde{f}_2(y, z, t),
\]
\[
\tilde{u}_x(x, 0, z, t) = \tilde{g}_1(x, z, t), \quad \tilde{u}_x(x, b, z, t) = \tilde{g}_2(x, z, t),
\]
\[
\tilde{u}_z(x, y, 0, t) = \tilde{h}_1(x, y, t), \quad \tilde{u}_z(x, y, c, t) = \tilde{h}_2(x, y, t),
\]
(3.3)

and the initial conditions
\[
\tilde{u}(x, y, z, 0) = \tilde{\varphi}(x, y, z). \quad (3.4)
\]

Second, we consider the fuzzy wave-like equation of the form
\[
\tilde{u}_{tt} = f(x, y, z) \circ \tilde{u}_{xx} \oplus g(x, y, z) \circ \tilde{u}_{yy} \oplus h(x, y, z) \circ \tilde{u}_{zz},
\]
\[
0 < x < a, 0 < y < b, 0 < z < c, t > 0, f(x, y, z) \geq 0, g(x, y, z) \geq 0, h(x, y, z) \geq 0,
\]
(3.5)

with the Neumann boundary conditions
\[
\tilde{u}_s(0, y, z, t) = \tilde{f}_1(y, z, t), \quad \tilde{u}_s(a, y, z, t) = \tilde{f}_2(y, z, t),
\]
\[
\tilde{u}_x(x, 0, z, t) = \tilde{g}_1(x, z, t), \quad \tilde{u}_x(x, b, z, t) = \tilde{g}_2(x, z, t),
\]
\[
\tilde{u}_z(x, y, 0, t) = \tilde{h}_1(x, y, t), \quad \tilde{u}_z(x, y, c, t) = \tilde{h}_2(x, y, t),
\]
(3.6)

and the initial conditions
\[
\tilde{u}(x, y, z, 0) = \tilde{\varphi}(x, y, z), \quad \tilde{u}_t(x, y, z, 0) = \tilde{\vartheta}(x, y, z).
\]
(3.7)

3.1 The fuzzy Adomian decomposition method
We derive the main algorithms of the fuzzy Adomian decomposition method for fuzzy heat-like and wave-like equations as follows.

3.1.1 Fuzzy heat-like equations
We consider the following parametric of fuzzy heat-like equation (3.1) of the form:
\[
\mu_t = f(x, y, z)\mu_{xx} + g(x, y, z)\mu_{yy} + h(x, y, z)\mu_{zz},
\]
\[
0 < x < a, 0 < y < b, 0 < z < c, t > 0, f(x, y, z) \geq 0, g(x, y, z) \geq 0, h(x, y, z) \geq 0,
\]
(3.8)

\[
\overline{\mu}_t = f(x, y, z)\overline{\mu}_{xx} + g(x, y, z)\overline{\mu}_{yy} + h(x, y, z)\overline{\mu}_{zz},
\]
\[
0 < x < a, 0 < y < b, 0 < z < c, t > 0, f(x, y, z) \geq 0, g(x, y, z) \geq 0, h(x, y, z) \geq 0.
\]
(3.9)

In an operator form, equations (3.8) and (3.9) can be written as
\[
L_t\mu = \left[f(x, y, z)\mu_{xx} + g(x, y, z)\mu_{yy} + h(x, y, z)\mu_{zz}\right](r),
\]
\[
0 < x < a, 0 < y < b, 0 < z < c, 0 < \alpha < 1, t > 0,
\]
(3.10)

\[
L_t\overline{\mu} = \left[f(x, y, z)\overline{\mu}_{xx} + g(x, y, z)\overline{\mu}_{yy} + h(x, y, z)\overline{\mu}_{zz}\right](r),
\]
\[
0 < x < a, 0 < y < b, 0 < z < c, 0 < \alpha < 1, t > 0,
\]
(3.11)
where the differential operator $L_t$ is

$$L_t = \frac{\partial}{\partial t}.$$  

Assuming that $L_t$ is invertible, the inverse operator $L_t^{-1}$ is given by the integral

$$L_t^{-1}(\cdot) = \int_0^t (\cdot) \, ds.$$  

(3.12)

Applying the inverse operator $L_t^{-1}$ to equations (3.10) and (3.11) and using the initial condition (3.4), we obtain

$$\begin{align*}
    u(x, y, z, t; r) &= \tilde{\phi}(x, y, z; r) + L_t^{-1} \left[ f(x, y, z) u_{xx} + g(x, y, z) u_{xy} + h(x, y, z) u_{xz} \right](r), \\
    \bar{u}(x, y, z, t; r) &= \tilde{\varphi}(x, y, z; r) + L_t^{-1} \left[ f(x, y, z) \bar{u}_{xx} + g(x, y, z) \bar{u}_{xy} + h(x, y, z) \bar{u}_{xz} \right](r).
\end{align*}$$

(3.13)

(3.14)

The fuzzy ADM consists of representing the solution for

$$\tilde{u}(x, y, z, t)(r) = \left[ u(x, y, z, t)(r), \bar{u}(x, y, z, t)(r) \right]$$

by an infinite sum of components as follows:

$$\begin{align*}
    u(x, y, z, t; r) &= \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r), \\
    \bar{u}(x, y, z, t; r) &= \sum_{m=0}^{\infty} \bar{u}_m(x, y, z, t; r),
\end{align*}$$

(3.15)

(3.16)

where the fuzzy components $\tilde{u}_m(x, y, z, t; r) = [\tilde{u}_m(x, y, z, t; r), \bar{u}_m(x, y, z, t; r)]$ will be recursively determined. Applying (3.16) to both sides of (3.13) gives

$$\begin{align*}
    \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) &= \tilde{\phi}(x, y, z; r) + L_t^{-1} \left[ \sum_{m=0}^{\infty} \left( \tilde{u}_m(x, y, z, t; r) \right) \right]_{xx} \\
    &+ g \left( \sum_{m=0}^{\infty} \left( \tilde{u}_m(x, y, z, t; r) \right) \right)_{xy} + h \left( \sum_{m=0}^{\infty} \left( \tilde{u}_m(x, y, z, t; r) \right) \right)_{xz},
\end{align*}$$

(3.17)

and

$$\begin{align*}
    \sum_{m=0}^{\infty} \bar{u}_m(x, y, z, t; r) &= \tilde{\varphi}(x, y, z; r) + L_t^{-1} \left[ \sum_{m=0}^{\infty} \left( \bar{u}_m(x, y, z, t; r) \right) \right]_{xx} \\
    &+ g \left( \sum_{m=0}^{\infty} \left( \bar{u}_m(x, y, z, t; r) \right) \right)_{xy} + h \left( \sum_{m=0}^{\infty} \left( \bar{u}_m(x, y, z, t; r) \right) \right)_{xz}.
\end{align*}$$

(3.18)

Next, the fuzzy ADM identifies the zeroth component $\tilde{u}_0(x, y, z, t; r) = [\tilde{u}_0(x, y, z, t; r), \bar{u}_0(x, y, z, t; r)]$ by all terms that arise from the initial condition and from integrating the source term if exists. To determine the components $\tilde{u}_m(x, y, z, t; r) = [\tilde{u}_m(x, y, z, t; r), \bar{u}_m(x, y, z, t; r)]$ by all terms that arise from the initial condition and from integrating the source term if exists.
\( z, t; r \), \( m \geq 1 \), of the solution \( \tilde{u}(x, y, z, t; r) \), fuzzy ADM formally introduces the recursive relation as follows:

\[
\begin{align*}
\tilde{u}_0(x, y, z, t; r) &= \varphi(x, y, z; r), \\
\tilde{u}_{m+1}(x, y, z, t; r) &= L_t^{-1}\left[ f(\tilde{u}_m(x, y, z, t; r))_{xx} + g(\tilde{u}_m(x, y, z, t; r))_{yy} + h(\tilde{u}_m(x, y, z, t; r))_{zz} \right], \quad m \geq 0,
\end{align*}
\]

(3.19)

Thus we can recurrently determine every component of the series \( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \) and so obtain the series solution of \( \tilde{u}(x, y, z, t; r) \).

### 3.1.2 Inhomogeneous fuzzy heat-like equations

We consider the inhomogeneous fuzzy heat-like equation of the form

\[
\tilde{u}_t = \tilde{\tau}(x, y, z) \oplus f(x, y, z) \odot \tilde{u}_{xx} \oplus g(x, y, z) \odot \tilde{u}_{yy} \oplus h(x, y, z) \odot \tilde{u}_{zz},
\]

\[0 < x < a, 0 < y < b, 0 < z < c, t > 0, f(x, y, z) \geq 0, g(x, y, z) \geq 0, h(x, y, z) \geq 0, \] (3.21)

The parametric form of (3.21) is

\[
L_t \tilde{u}(r) = \tilde{\tau}(x, y, z; r) + \left[ f(x, y, z) u_{xx} + g(x, y, z) u_{yy} + h(x, y, z) u_{zz} \right](r),
\]

\[0 < x < a, 0 < y < b, 0 < z < c, 0 < \alpha < 1, t > 0, \] (3.22)

\[
L_t \tilde{u}(r) = \tilde{\tau}(x, y, z; r) + \left[ f(x, y, z) \tilde{u}_{xx} + g(x, y, z) \tilde{u}_{yy} + h(x, y, z) \tilde{u}_{zz} \right](r),
\]

\[0 < x < a, 0 < y < b, 0 < z < c, 0 < \alpha < 1, t > 0, \] (3.23)

where

\[
L_t = \frac{\partial}{\partial t}.
\]

Assuming that \( L_t \) is invertible, the inverse operator \( L_t^{-1} \) is given by the integral

\[
L_t^{-1}(\cdot) = \int_0^t (\cdot) ds.
\]

(3.24)

Applying the inverse operator \( L_t^{-1} \) to (3.22) and (3.23), by the initial condition (3.4) we obtain

\[
\tilde{u}(x, y, z, t; r) = \tilde{\tau}(x, y, z; r) + L_t^{-1}\left[ f(x, y, z) \tilde{u}_{xx} + g(x, y, z) \tilde{u}_{yy} + h(x, y, z) \tilde{u}_{zz} \right](r),
\]

(3.25)
\[ \Pi(x, y, z, t; r) = \tau(x, y, z; r)t + L_i^{-1} \left[ f(x, y, z) \Pi_{xx} + g(x, y, z) \Pi_{yy} + h(x, y, z) \Pi_{zz} \right] (r). \] (3.26)

The fuzzy ADM assumes a series solution for
\[ \tilde{u}(x, y, z, t; r) = \left[ u(x, y, z, t; r), \Pi(x, y, z, t; r) \right], \]
given by an infinite sum of components as follows:
\[ u(x, y, z, t; r) = \sum_{m=0}^{\infty} u_m(x, y, z, t; r), \] (3.27)
\[ \Pi(x, y, z, t; r) = \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r), \] (3.28)
where the fuzzy components \( \tilde{u}_m(x, y, z, t; r) = [u_m(x, y, z, t; r), \Pi_m(x, y, z, t; r)] \) will be recursively determined. Substituting (3.28) into both sides of (3.25) gives
\[
\sum_{m=0}^{\infty} u_m(x, y, z, t; r) = \tau(x, y, z; r)t + L_i^{-1} \left[ f \left( \sum_{m=0}^{\infty} u_m(x, y, z, t; r) \right) \right]_{xx} + g \left( \sum_{m=0}^{\infty} u_m(x, y, z, t; r) \right)_{yy} + h \left( \sum_{m=0}^{\infty} u_m(x, y, z, t; r) \right)_{zz} \] (3.29)

and
\[
\sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) = \tau(x, y, z; r)t + L_i^{-1} \left[ f \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) \right) \right]_{xx} + g \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) \right)_{yy} + h \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) \right)_{zz}. \] (3.30)

Following the fuzzy ADM, we identify the zeroth component \( \tilde{u}_0(x, y, z, t; r) = [u_0(x, y, z, t; r), \Pi_0(x, y, z, t; r)] \) by all terms that arise from the initial condition and from integrating the source term if exists. To determine the components \( \tilde{u}_m(x, y, z, t; r), m \geq 1, \) of the solution \( \tilde{u}(x, y, z, t; r), \) the fuzzy ADM formally introduces the recursive relation
\[
u_0(x, y, z, t; r) = \tau(x, y, z; r)t, \]
\[
u_{m+1}(x, y, z, t) = L_i^{-1} \left[ f(u_m(x, y, z, t; r)) \right]_{xx} + g(u_m(x, y, z, t; r))_{yy} + h(u_m(x, y, z, t; r))_{zz}, \] (3.31)
\[ m \geq 0, \]
\[
u_0(x, y, z, t; r) = \tau(x, y, z; r)t, \]
\[
u_{m+1}(x, y, z, t) = L_i^{-1} \left[ f(\Pi_m(x, y, z, t; r)) \right]_{xx} + g(\Pi_m(x, y, z, t; r))_{yy} + h(\Pi_m(x, y, z, t; r))_{zz}, \] (3.32)
\[ m \geq 0. \]
3.1.3 Fuzzy wave-like equations

We consider the fuzzy wave-like equation (3.5) of the form

\[ u_{tt} = [f(x, y, z)u_{xx} + g(x, y, z)u_{yy} + h(x, y, z)u_{zz}] (r), \]

\[ 0 < x < a, 0 < y < b, 0 < z < c, 1 < \alpha \leq 2, t > 0, \] (3.33)

\[ \tilde{u}_{tt} = [f(x, y, z)\tilde{u}_{xx} + g(x, y, z)\tilde{u}_{yy} + h(x, y, z)\tilde{u}_{zz}] (r), \]

\[ 0 < x < a, 0 < y < b, 0 < z < c, 1 < \alpha \leq 2, t > 0. \] (3.34)

In an operator form, equations (3.33) and (3.34) become

\[ L_t u(r) = [f(x, y, z)u_{xx} + g(x, y, z)u_{yy} + h(x, y, z)u_{zz}] (r), \] (3.35)

\[ L_t \tilde{u}(r) = [f(x, y, z)\tilde{u}_{xx} + g(x, y, z)\tilde{u}_{yy} + h(x, y, z)\tilde{u}_{zz}] (r), \] (3.36)

where

\[ L_t = \frac{\partial^2}{\partial t^2}. \]

Assuming that \( L_t \) is invertible, the inverse operator \( L_t^{-1} \) is given by the integral

\[ L_t^{-1} (\cdot) = \int_0^t \int_0^t (\cdot) ds \, ds. \] (3.37)

Applying the inverse operator \( L_t^{-1} \) to both sides of equations (3.35) and (3.36), by the initial conditions of equation (3.7) we have

\[ u(x, y, z, t; r) = \psi(x, y, z; r) + t \tilde{\psi}(x, y, z; r) \]

\[ + L_t^{-1} [f(x, y, z)u_{xx} + g(x, y, z)u_{yy} + h(x, y, z)u_{zz}] (r), \] (3.38)

\[ \tilde{u}(x, y, z, t; r) = \varphi(x, y, z; r) + t \tilde{\varphi}(x, y, z; r) \]

\[ + L_t^{-1} [f(x, y, z)\tilde{u}_{xx} + g(x, y, z)\tilde{u}_{yy} + h(x, y, z)\tilde{u}_{zz}] (r). \] (3.39)

The fuzzy ADM assumes a series solution for

\[ \tilde{u}(x, y, z; t; r) = \left[ u(x, y, z; t; r), \tilde{u}(x, y, z; t; r) \right], \]

given by an infinite sum of components as follows:

\[ u(x, y, z, t; r) = \sum_{m=0}^{\infty} u_m(x, y, z; t; r), \] (3.40)

\[ \tilde{u}(x, y, z, t; r) = \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z; t; r), \] (3.41)
where the fuzzy components $\tilde{u}_m(x, y, z, t; r) = [\tilde{u}_m(x, y, z, t; r), \tilde{u}_m(x, y, z, t; r)]$ will be repeatedly identified. Applying (3.41) to both sides of (3.38) gives

$$
\sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) = \varphi(x, y, z; r) + t\tilde{\varphi}(x, y, z; t; r) + L_t^{-1}\left[ f\left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \right) \right]_{xx} + g\left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \right) + h\left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \right)_{yy} \right] (3.42)
$$

and

$$
\sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) = \varphi(x, y, z; r) + t\tilde{\varphi}(x, y, z; t; r) + L_t^{-1}\left[ f\left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \right) \right]_{xx} + g\left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \right) + h\left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, z, t; r) \right)_{yy} \right] (3.43)
$$

To specify the components $\tilde{u}_m(x, y, z, t; r)$, $m \geq 0$, of the solution $\tilde{u}(x, y, z; r)$, the fuzzy ADM officially introduces the recursive relationship

$$
u \mu(x, y, z, t; r) = \varphi(x, y, z; r) + t\tilde{\varphi}(x, y, z; t; r),
$$

$$
\tilde{u}_{m+1}(x, y, z, t; r) = L_t^{-1}\left[ f\left( \tilde{u}_m(x, y, z, t; r) \right) \right]_{xx} + g\left( \tilde{u}_m(x, y, z, t; r) \right) + h\left( \tilde{u}_m(x, y, z, t; r) \right)_{yy} \right] (3.44)
$$

$$
\tilde{u}_0(x, y, z, t; r) = \varphi(x, y, z; r) + t\tilde{\varphi}(x, y, z; t; r),
$$

$$
\tilde{u}_{m+1}(x, y, z, t; r) = L_t^{-1}\left[ f\left( \tilde{u}_m(x, y, z, t; r) \right) \right]_{xx} + g\left( \tilde{u}_m(x, y, z, t; r) \right) + h\left( \tilde{u}_m(x, y, z, t; r) \right)_{yy} \right] (3.45)
$$

With this relationship, we easily obtain the components $\tilde{u}_m(x, y, z, t; r)$, $m \geq 0$. This, in turn, leads to a solution in the form of a series, but the solution is in a closed-form, which directly follows if there is an exact solution.

### 3.2 The fuzzy variational iteration method

We derive the main algorithms of the fuzzy variational iteration method considering the following general fuzzy partial differential equation:

$$
L_t\tilde{u} \oplus L_x\tilde{u} \oplus L_y\tilde{u} \oplus L_z\tilde{u} \oplus N\tilde{u} = \tilde{g}(x, y, z, t),
$$

where $L_t, L_x, L_y$, and $L_z$ are linear operators of $t, x, y$, and $z$, respectively, and $N$ is a non-linear operator. Applying the fuzzy variational iteration method, we can express the following
correction functional (in \( t, \ x, \ y, \) and \( z \)-directions), respectively, as

\[
\begin{align*}
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^t \lambda_1 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds, \\
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^x \lambda_2 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds, \\
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^y \lambda_3 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds, \\
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^z \lambda_4 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds,
\end{align*}
\]

(3.47)

and

\[
\begin{align*}
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^t \lambda_1 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds, \\
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^x \lambda_2 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds, \\
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^y \lambda_3 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds, \\
\mathcal{U}_{m+1}(x, y, z, t; r) &= \mathcal{U}_m(x, y, z, t; r) + \int_0^z \lambda_4 [L_x \mathcal{U}_m + (L_x + L_y + L_z + N) \mathcal{U}_m - g] \, ds,
\end{align*}
\]

(3.48)

where \( \lambda_1, \lambda_2, \lambda_3, \) and \( \lambda_4 \) are general Lagrange multipliers, which can specified optimally via the fuzzy variational theory, and

\[
\mathcal{U}_m(x, y, z, t; r) = [\mathcal{U}_m(x, y, z, t; r), \mathcal{U}_m(x, y, z, t; r)] \text{ via restricted variation with } \delta \mathcal{U}_m = 0 \text{ and } \delta \mathcal{U}_m = 0.\]

Using the fuzzy variational iteration method, we first determine the Lagrange multipliers \( \lambda_i, \ i = 1, 2, 3, 4. \) We will easily obtain sequential approximations \( \delta \mathcal{U}_{m+1}, m \geq 0, \) from solution \( \mathcal{U}_m(x, y, z, t; r) = [u(x, y, z, t; r), \mathcal{U}_m(x, y, z, t; r)] \), that is, a suitable choice for an experimental function that has an ambiguous value

\[
\mathcal{U}_m(x, y, z, t; r) = [u_m(x, y, z, t; r), \mathcal{U}_m(x, y, z, t; r)].
\]

Then the solution is given as follows:

\[
\begin{align*}
u(x, y, z, t; r) &= \lim_{m \to \infty} \mathcal{U}_m(x, y, z, t; r), \\
\mathcal{n}(x, y, z, t; r) &= \lim_{m \to \infty} \mathcal{U}_m(x, y, z, t; r).
\end{align*}
\]

(3.49)

(3.50)

3.2.1 Fuzzy heat-like equations

We consider the fuzzy heat-like equation of the form

\[
\tilde{u}_t = g_1(x, y, z) \odot \tilde{u}_x \odot g_2(x, y, z) \odot \tilde{u}_y \odot g_3(x, y, z) \odot \tilde{u}_z
\]

(3.51)

subject to the initial condition

\[
\tilde{u}(x, y, z, 0) = \tilde{g}_4(x, y, z).
\]

(3.52)

The parametric form of (3.51) is

\[
\tilde{u}_t = g_1(x, y, z) \tilde{u}_x + g_2(x, y, z) \tilde{u}_y + g_3(x, y, z) \tilde{u}_z,
\]

(3.53)
\[ u_t = g_1(x, y, z)u_{xx} + g_2(x, y, z)u_{yy} + g_3(x, y, z)u_{zz} \]  

(3.54)

We construct a correction functional for equations (3.53) and (3.54) applying the fuzzy variational iteration method of the form

\[ u_{m+1}(x, y, z; t; r) = u_m(x, y, z; t; r) + \int_0^t \lambda [ (u_m)_t + g_1(x, y, z)(u_m)_{xx} + g_2(x, y, z)(u_m)_{yy} + g_3(x, y, z)(u_m)_{zz} ] ds, \]  

(3.55)

\[ \bar{u}_{m+1}(x, y, z; t; r) = \bar{u}_m(x, y, z; t; r) + \int_0^t \lambda [ (\bar{u}_m)_{t} + g_1(x, y, z)(\bar{u}_m)_{xx} + g_2(x, y, z)(\bar{u}_m)_{yy} + g_3(x, y, z)(\bar{u}_m)_{zz} ] ds, \]  

(3.56)

where \( m \geq 0 \), and

\[ u_0(x, y, z; t; r) = u(x, y, z; 0; r), \]

\[ \bar{u}_0(x, y, z; t; r) = \bar{u}(x, y, z; 0; r). \]

The fuzzy variational iteration method solution \( \tilde{u}(x, y, z; t; r) = [u(x, y, z; t; r), \bar{u}(x, y, z; t; r)] \) takes the form

\[ \tilde{u}(x, y, z; t; r) = \lim_{m \to \infty} u_m(x, y, z; t; r), \]

\[ \bar{u}(x, y, z; t; r) = \lim_{m \to \infty} \bar{u}_m(x, y, z; t; r), \]

where \( \tilde{u}_m(x, y, z; t; r) = [u_m(x, y, z; t; r), \bar{u}_m(x, y, z; t; r)] \) will be recursively specified.

3.2.2 Inhomogeneous fuzzy heat-like equations

We consider the inhomogeneous fuzzy heat-like equation of the form

\[ \tilde{u}_t = \tilde{\xi}(x, y, z) \oplus g_1(x, y, z) \odot \tilde{u}_{xx} \oplus g_2(x, y, z) \odot \tilde{u}_{yy} \oplus g_3(x, y, z) \odot \tilde{u}_{zz} \]  

(3.57)

subject to the initial condition

\[ \tilde{u}(x, y, z; 0) = \tilde{g}_4(x, y, z). \]  

(3.58)

The parametric form of (3.57) is

\[ \tilde{u}_t = \tilde{\xi}(x, y, z) + g_1(x, y, z)u_{xx} + g_2(x, y, z)u_{yy} + g_3(x, y, z)u_{zz}, \]  

(3.59)

\[ \bar{u}_t = \bar{\xi}(x, y, z) + g_1(x, y, z)\bar{u}_{xx} + g_2(x, y, z)\bar{u}_{yy} + g_3(x, y, z)\bar{u}_{zz}. \]  

(3.60)

We construct a correction functional for equations (3.59) and (3.60) applying the variational iteration method of the form

\[ u_{m+1}(x, y, z; t; r) = u_m(x, y, z; t; r) + \int_0^t \lambda [ (u_m(x, y, z; t; r)) - \tilde{\xi}(x, y, z) - F ] ds, \]  

(3.61)
\[ \bar{u}_{m+1}(x, y, z, t; r) = \bar{u}_m(x, y, z, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, y, z, t; r) \right)_t - \overline{f}(x, y, z) - \overline{F} \right] ds, \]  
\hspace{1cm} (3.62)

where \( m \geq 0, \)

\[ \bar{u}_0(x, y, z, t; r) = u(x, y, z, 0; r), \]
\[ \overline{u}_0(x, y, z, t; r) = \overline{u}(x, y, z, 0; r), \]

and the above-mentioned \( \xi(x, y, z; r) = [\xi(x, y, z; r), \overline{\xi}(x, y, z; r)] \) is the fuzzy inhomogeneous term; however, \( \overline{F} \) and \( \overline{F} \) are

\[ \overline{F} = g_1(\bar{u}_m(x, y, z, t; r))_{xx} + g_2(\bar{u}_m(x, y, z, t; r))_{yy} + g_3(\bar{u}_m(x, y, z, t; r))_{zz}, \]  
\hspace{1cm} (3.63)

\[ \overline{F} = g_1(\overline{u}_m(x, y, z, t; r))_{xx} + g_2(\overline{u}_m(x, y, z, t; r))_{yy} + g_3(\overline{u}_m(x, y, z, t; r))_{zz}. \]  
\hspace{1cm} (3.64)

The fuzzy variational iteration method solution \( \tilde{u}(x, y, z, t; r) = [u(x, y, z, t; r), \overline{u}(x, y, z, t; r)] \) takes the form

\[ \tilde{u}(x, y, z, t; r) = \lim_{m \to \infty} \bar{u}_m(x, y, z, t; r), \]
\[ \overline{u}(x, y, z, t; r) = \lim_{m \to \infty} \overline{u}_m(x, y, z, t; r), \]

where \( \tilde{u}_m(x, y, z, t; r) = [\bar{u}_m(x, y, z, t; r), \overline{u}_m(x, y, z, t; r)] \) will be recursively specified.

3.2.3 Fuzzy wave-like equations

We consider the fuzzy wave-like equation of the form

\[ \tilde{u}_{tt} = h_1(x, y, z) \odot \tilde{u}_{xx} + h_2(x, y, z) \odot \tilde{u}_{yy} + h_3(x, y, z) \odot \tilde{u}_{zz} \]  
\hspace{1cm} (3.65)

subject to the initial condition

\[ \tilde{u}(x, y, z, 0) = \tilde{h}_4(x, y, z), \hspace{0.5cm} \tilde{u}_t(x, y, z, 0) = \tilde{h}_5(x, y, z). \]  
\hspace{1cm} (3.66)

The parametric form of (3.65) is

\[ \bar{u}_t = h_1(x, y, z) \bar{u}_{xx} + h_2(x, y, z) \bar{u}_{yy} + h_3(x, y, z) \bar{u}_{zz}, \]  
\hspace{1cm} (3.67)

\[ \overline{u}_t = h_1(x, y, z) \overline{u}_{xx} + h_2(x, y, z) \overline{u}_{yy} + h_3(x, y, z) \overline{u}_{zz}. \]  
\hspace{1cm} (3.68)

According to the variational iteration method, equations (3.67) and (3.68) take the form

\[ \bar{u}_{m+1}(x, y, z, t; r) = \bar{u}_m(x, y, z, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, y, z, t; r) \right)_t + h_1(x, y, z)(\bar{u}_m)_{xx} + h_2(x, y, z)(\bar{u}_m)_{yy} + h_3(x, y, z)(\bar{u}_m)_{zz} \right] ds, \]  
\hspace{1cm} (3.69)

\[ \overline{u}_{m+1}(x, y, z, t; r) = \overline{u}_m(x, y, z, t; r) + \int_0^t \lambda \left[ \left( \overline{u}_m(x, y, z, t; r) \right)_t + h_1(x, y, z)(\overline{u}_m)_{xx} + h_2(x, y, z)(\overline{u}_m)_{yy} + h_3(x, y, z)(\overline{u}_m)_{zz} \right] ds, \]  
\hspace{1cm} (3.70)
where \( m \geq 0 \), and \( \tilde{u}_0(x,y,z,t;r) = [u(x,y,z,0;r) + tu_u(x,y,z,0;r), \Pi(x,y,z,0;r) + t\Pi_t(x,y,z,0;r)] \).

The fuzzy variational iteration method solution \( \tilde{u}(x,y,z,t;r) = [u(x,y,z,t;r), \Pi(x,y,z,t;r)] \) can be expressed as

\[
\tilde{u}(x,y,z,t;r) = \lim_{m \to \infty} \tilde{u}_m(x,y,z,t;r),
\]

\[
\Pi(x,y,z,t;r) = \lim_{m \to \infty} \Pi_m(x,y,z,t;r),
\]

where \( \tilde{u}_m(x,y,z,t;r) = [u_m(x,y,z,t;r), \Pi_m(x,y,z,t;r)] \) will be recursively specified.

### 4 Examples

In this section, we demonstrate how fuzzy Adomian decomposition method and fuzzy variational iteration method can be easily applied to get the exact solutions of the fuzzy heat-like and wave-like equations with variable coefficients.

**Example 4.1** We consider the following one-dimensional fuzzy heat-like equation with variable coefficients:

\[
\tilde{u}_t = \frac{1}{2} x^2 \otimes \tilde{u}_{xx}, \quad 0 < x < 1, t > 0,
\]  

(4.1)

with the boundary conditions

\[
\tilde{u}(0,t) = 0, \quad \tilde{u}(1,t) = \exp(t)
\]

(4.2)

and the initial condition

\[
\tilde{u}(x,0) = \left[(-0.5 + 1.5r)^n, (1.5 - 0.5r)^n\right] \otimes x^2,
\]

(4.3)

where \( n = 1, 2, 3, \ldots \).

The parametric form of (4.1) is

\[
\tilde{u}_t = \frac{1}{2} x^2 \tilde{u}_{xx}, \quad 0 < x < 1, t > 0,
\]

(4.4)

\[
\Pi_t = \frac{1}{2} x^2 \Pi_{xx}, \quad 0 < x < 1, t > 0,
\]

(4.5)

for \( r \in [0,1] \), where \( u \) stands for \( u(x,t;r) \), and similarly \( \Pi \).

**Case [A]. Fuzzy Adomian decomposition method**

In an operator form, equations (4.4) and (4.5) are

\[
L_u \tilde{u} = \frac{1}{2} x^2 \tilde{u}_{xx}, \quad 0 < x < 1,
\]

(4.6)

\[
L_\Pi \Pi = \frac{1}{2} x^2 \Pi_{xx}, \quad 0 < x < 1.
\]

(4.7)

Using the initial condition (4.3) and applying \( L_u^{-1} \) to equations (4.6) and (4.7), we get

\[
u(x,t;r) = (1.5r - 0.5)^n x^2 + L_u^{-1} \left[ \frac{1}{2} x^2 \tilde{u}_{xx} \right],
\]

(4.8)
\[ u(x,t;r) = (1.5 - 0.5r)x^2 + L_t^{-1} \left[ \frac{1}{2} x^2 u_{xx} \right]. \] (4.9)

Substituting the ADM series (3.15) and (3.16) for \( \tilde{u} \) into (4.8) and (4.9), we define
\[
\sum_{m=0}^{\infty} u_m(x,t;r) = (1.5r - 0.5)x^2 + L_t^{-1} \left[ \frac{1}{2} x^2 \left( \sum_{m=0}^{\infty} u_m(x,t;r) \right) \right] \] , (4.10)
\[
\sum_{m=0}^{\infty} \bar{u}_m(x,t;r) = (1.5 - 0.5r)x^2 + L_t^{-1} \left[ \frac{1}{2} x^2 \left( \sum_{m=0}^{\infty} \bar{u}_m(x,t;r) \right) \right] . \] (4.11)

Next, we use the recursive relations
\[
u_0(x,t;r) = (1.5r - 0.5)x^2, \]
\[ u_{m+1}(x,t;r) = L_t^{-1} \left[ \frac{1}{2} x^2 (u_m)_{xx} \right], \quad j \geq 0 \] \] (4.12)
and
\[
\bar{u}_0(x,t;r) = (1.5 - 0.5r)x^2, \]
\[ \bar{u}_{m+1}(x,t;r) = L_t^{-1} \left[ \frac{1}{2} x^2 (\bar{u}_m)_{xx} \right], \quad j \geq 0 \] \] (4.13)

This in turn gives the first few components:
\[
u_0(x,t;r) = (1.5r - 0.5)x^2, \]
\[ u_1(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (u_0)_{xx} \right] = (1.5r - 0.5)x^2 t, \]
\[ u_2(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (u_1)_{xx} \right] = (1.5r - 0.5)x^2 t^2 \]
\[ u_3(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (u_2)_{xx} \right] = (1.5r - 0.5)x^2 \frac{t^3}{3!} \]
\[ \vdots \]
\[ u_m(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (u_{m-1})_{xx} \right] = (1.5r - 0.5)x^2 \frac{t^m}{m!} \]
and
\[
\bar{u}_0(x,t;r) = (1.5 - 0.5r)x^2, \]
\[ \bar{u}_1(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (\bar{u}_0)_{xx} \right] = (1.5 - 0.5r)x^2 t, \]
\[ \bar{u}_2(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (\bar{u}_1)_{xx} \right] = (1.5 - 0.5r)x^2 \frac{t^2}{2!} \]
\[ \bar{u}_3(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (\bar{u}_2)_{xx} \right] = (1.5 - 0.5r)x^2 \frac{t^3}{3!} \]
\[ \vdots \]
\[ \bar{u}_m(x,t;r) = L_t^{-1} \left[ \frac{1}{2} L_t^{-1} x^2 (\bar{u}_{m-1})_{xx} \right] = (1.5 - 0.5r)x^2 \frac{t^m}{m!} \]
Consequently, the solution for the standard fuzzy heat-like equation for $\alpha = 1$ is given by

$$u(x, t; r) = (1.5r - 0.5)^n \left[ x^2 \left( 1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \ldots + \frac{t^n}{n!} \right) \right],$$

$$\bar{u}(x, t; r) = (1.5 - 0.5r)^n \left[ x^2 \left( 1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \ldots + \frac{t^n}{n!} \right) \right].$$

and in the closed form by

$$\tilde{u}(x, t; r) = \left[ (1.5r - 0.5)^n, (1.5 - 0.5r)^n \right] \odot (\exp(t)x^2), \quad 0 \leq r \leq 1.$$

Case [B]. Fuzzy variational iteration method

The fuzzy VIM schema of equations (4.4) and (4.5) are

$$u_{m+1}(x, t; r) = u_m(x, t; r) + \int_0^t \lambda \left[ \left( u_m(x, s; r) \right)_s - \frac{x^2}{2} \left( u_m(x, s; r) \right)_{xx} \right] ds, \quad (4.14)$$

$$\bar{u}_{m+1}(x, t; r) = \bar{u}_m(x, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, s; r) \right)_s - \frac{x^2}{2} \left( \bar{u}_m(x, s; r) \right)_{xx} \right] ds. \quad (4.15)$$

Making the above correction functional stationary, we have

$$\delta u_{m+1}(x, t; r) = \delta u_m(x, t; r) + \int_0^t \lambda \left[ \left( u_m(x, s; r) \right)_s - \frac{x^2}{2} \left( u_m(x, s; r) \right)_{xx} \right] ds, \quad (4.16)$$

$$\delta \bar{u}_{m+1}(x, t; r) = \delta \bar{u}_m(x, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, s; r) \right)_s - \frac{x^2}{2} \left( \bar{u}_m(x, s; r) \right)_{xx} \right] ds. \quad (4.17)$$

Therefore the stationary conditions are given by

$$1 + \lambda|_{s=2} = 0, \quad \lambda'(s) = 0.$$

By integrating the differential equation we have

$$\lambda'(s) = 0 \quad \Rightarrow \quad \lambda(s) = c_1.$$

Using the condition, we arrive at

$$1 + \lambda|_{s=2} = 0 \quad \Rightarrow \quad c_1 = -1.$$

Therefore the Lagrange multiplier $\lambda$ is given by

$$\lambda = -1.$$

Substituting the value of the Lagrange multiplier into equations (4.14) and (4.15) gives the following iterative formulas:

$$u_{m+1}(x, t; r) = u_m(x, t; r) - \int_0^t \left[ \left( u_m(x, s; r) \right)_s - \frac{x^2}{2} \left( u_m(x, s; r) \right)_{xx} \right] ds, \quad (4.18)$$
\[ \begin{aligned}
\pi_{m+1}(x, t; r) &= \pi_m(x, t; r) - \int_0^t \left[ \left( \frac{x^2}{2} \right) (\pi_m(x, s; r))_{xx} \right] ds. 
\end{aligned} \tag{4.19} \]

Using iteration formulas (4.18) and (4.19) with the initial approximation \( \tilde{u}_0(x, t; r) = \tilde{u}(x, 0; r) = [(-0.5 + 1.5r)^n, (1.5 - 0.5r)^n] \odot x^2 \), we get the following successive approximations:

\[ \begin{aligned}
\tilde{u}_1(x, t; r) &= (1.5r - 0.5)^n(1 + t)x^2, \\
\tilde{u}_2(x, t; r) &= (1.5r - 0.5)^n\left(1 + t + \frac{t^2}{2!}\right)x^2, \\
\tilde{u}_3(x, t; r) &= (1.5r - 0.5)^n\left(1 + t + \frac{t^2}{2!} + \frac{t^3}{3!}\right)x^2, \\
&\vdots \\
\tilde{u}_m(x, t; r) &= (1.5r - 0.5)^n\left(1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots + \frac{t^m}{m!}\right)x^2, 
\end{aligned} \tag{4.20} \]

and

\[ \begin{aligned}
\bar{u}_1(x, t; r) &= (1.5 - 0.5r)^n(1 + t)x^2, \\
\bar{u}_2(x, t; r) &= (1.5 - 0.5r)^n\left(1 + t + \frac{t^2}{2!}\right)x^2, \\
\bar{u}_3(x, t; r) &= (1.5 - 0.5r)^n\left(1 + t + \frac{t^2}{2!} + \frac{t^3}{3!}\right)x^2, \\
&\vdots \\
\bar{u}_m(x, t; r) &= (1.5 - 0.5r)^n\left(1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots + \frac{t^m}{m!}\right)x^2, 
\end{aligned} \tag{4.21} \]

Consequently, the exact solution is given by

\[ \begin{aligned}
\bar{u}(x, t; r) &= (1.5r - 0.5)^n \lim_{m \to \infty} \left[ x^2 \left(1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots + \frac{t^m}{m!}\right) \right], \\
\underline{u}(x, t; r) &= (1.5 - 0.5r)^n \lim_{m \to \infty} \left[ x^2 \left(1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots + \frac{t^m}{m!}\right) \right], 
\end{aligned} \]

and we can rewrite the solution as follows:

\[ \begin{aligned}
\tilde{u}(x, t; r) &= \left[(1.5r - 0.5)^n, (1.5 - 0.5r)^n\right] \odot (\exp(t)x^2), \quad 0 \leq r \leq 1.
\end{aligned} \]

\textbf{Example 4.2} We consider the following two-dimensional fuzzy heat-like equation with variable coefficients:

\[ \tilde{u}_t = \frac{1}{2} \left( y^2 \odot \tilde{u}_{yy} + x^2 \odot \tilde{u}_{xx} \right), \quad 0 < x, y < 1, t > 0, \tag{4.22} \]
with the Neumann boundary conditions
\[ \tilde{u}_x(0, y, t) = 0, \quad \tilde{u}_x(1, y, t) = 2 \sinh t, \]
\[ \tilde{u}_y(x, 0, t) = 0, \quad \tilde{u}_y(x, 1, t) = 2 \coth t, \]  
and the initial condition
\[ \tilde{u}(x, y, 0) = \left[(1 + 2r)^n, (5 - 2r^n)\right] \otimes g(t) y^2, \]  
where \( n = 1, 2, 3, \ldots \).

The parametric form of (4.22) is
\[ u_t = \frac{1}{2} \left(y^2 u_{xx} + x^2 u_{yy}\right), \quad 0 < x, y < 1, t > 0, \]  
\[ \bar{u}_t = \frac{1}{2} \left(y^2 \bar{u}_{xx} + x^2 \bar{u}_{yy}\right), \quad 0 < x, y < 1, t > 0, \]  
for \( r \in [0, 1] \), where \( \mu \) stands for \( u(x, y, t; r) \), and similarly for \( \bar{u} \).

Case [A]. **Fuzzy Adomian decomposition method**

In an operator form, equations (4.25) and (4.26) are
\[ L_t \mu = \frac{1}{2} \left(y^2 \mu_{xx} + x^2 \mu_{yy}\right), \]  
\[ L_t \bar{u} = \frac{1}{2} \left(y^2 \bar{u}_{xx} + x^2 \bar{u}_{yy}\right). \]  
Applying the inverse operator \( L^{-1} \) to both sides of equations (4.27) and (4.28), we obtain
\[ \mu(x, y, t; r) = (1 + 2r)^n - y^2 + \frac{1}{2} L^{-1}_t \left(y^2 \mu_{xx} + x^2 \mu_{yy}\right), \]  
\[ \bar{u}(x, y, t; r) = (5 - 2r^n) - y^2 + \frac{1}{2} L^{-1}_t \left(y^2 \bar{u}_{xx} + x^2 \bar{u}_{yy}\right). \]  
Using the fuzzy ADM series for \( \tilde{u}(x, y, t; r) = [\mu(x, y, t; r), \bar{u}(x, y, t; r)] \), we get
\[ \sum_{m=0}^{\infty} \mu_m(x, y, t; r) = (1 + 2r)^n - y^2 + \frac{1}{2} L^{-1}_t \left[y^2 \left(\sum_{m=0}^{\infty} \mu_m(x, y, t; r)\right)_{xx}\right] \]  
\[ + x^2 \left(\sum_{m=0}^{\infty} \mu_m(x, y, t; r)\right)_{yy}, \]  
\[ \sum_{m=0}^{\infty} \bar{u}_m(x, y, t; r) = (5 - 2r^n) - y^2 + \frac{1}{2} L^{-1}_t \left[y^2 \left(\sum_{m=0}^{\infty} \bar{u}_m(x, y, t; r)\right)_{xx}\right] \]  
\[ + x^2 \left(\sum_{m=0}^{\infty} \bar{u}_m(x, y, t; r)\right)_{yy}, \]  
where \( m \geq 0 \).
Now we use the recursive relations

\[
\begin{aligned}
\bar{u}_0(x, y; t; r) &= (1 + 2r)^n - y^2, \\
\bar{u}_m(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left( y^2 (\bar{u}_m)_{xx} + x^2 (\bar{u}_m)_{yy} \right), \quad m \geq 0,
\end{aligned}
\] (4.33)

\[
\begin{aligned}
\tilde{u}_0(x, y; t; r) &= (5 - 2r)^n - y^2, \\
\tilde{u}_m(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left( y^2 (\tilde{u}_m)_{xx} + x^2 (\tilde{u}_m)_{yy} \right), \quad m \geq 0,
\end{aligned}
\] (4.34)

and thus the first few components are

\[
\begin{aligned}
\bar{u}_0(x, y; t; r) &= (1 + 2r)^n - y^2, \\
\bar{u}_1(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left[ y^2 (\bar{u}_0)_{xx} + x^2 (\bar{u}_0)_{yy} \right] = (1 + 2r)^n - x^2 t, \\
\bar{u}_2(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left[ y^2 (\bar{u}_1)_{xx} + x^2 (\bar{u}_1)_{yy} \right] = (1 + 2r)^n - y^2 \frac{t^2}{2!}, \\
\bar{u}_3(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left[ y^2 (\bar{u}_2)_{xx} + x^2 (\bar{u}_2)_{yy} \right] = (1 + 2r)^n - x^2 \frac{t^3}{3!}, \\
&\vdots
\end{aligned}
\]

and

\[
\begin{aligned}
\tilde{u}_0(x, y; t; r) &= (5 - 2r)^n - y^2, \\
\tilde{u}_1(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left[ y^2 (\tilde{u}_0)_{xx} + x^2 (\tilde{u}_0)_{yy} \right] = (5 - 2r)^n - x^2 t, \\
\tilde{u}_2(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left[ y^2 (\tilde{u}_1)_{xx} + x^2 (\tilde{u}_1)_{yy} \right] = (5 - 2r)^n - y^2 \frac{t^2}{2!}, \\
\tilde{u}_3(x, y; t; r) &= \frac{1}{2} L^{-1}_t \left[ y^2 (\tilde{u}_2)_{xx} + x^2 (\tilde{u}_2)_{yy} \right] = (5 - 2r)^n - x^2 \frac{t^3}{3!}, \\
&\vdots
\end{aligned}
\]

Thus the solution in a series form is

\[
\begin{aligned}
\bar{u}(x, y; t; r) &= (1 + 2r)^n - \left[ x^2 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots \right) + y^2 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots \right) \right], \\
\tilde{u}(x, y; t; r) &= (5 - 2r)^n - \left[ x^2 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots \right) + y^2 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots \right) \right],
\end{aligned}
\]

and then the exact solution is

\[
\bar{u}(x, y; t; r) = \left[ (1 + 2r)^n, (5 - 2r)^n \right] \otimes_{GH} (x^2 \sinh t + y^2 \cosh t), \quad 0 \leq r \leq 1.
\]

Case [B]. Fuzzy variational iteration method
The correction functionals for equations (4.25) and (4.26) are

\[ \mathcal{U}_{m+1}(x, y; t; r) = \mathcal{U}_m(x, y; t; r) + \int_0^t \lambda \left( \left( \mathcal{U}_m(x, y; s; r) \right)_s - \frac{\gamma^2}{2} \left( \mathcal{U}_m(x, y; s; r) \right)_{xx} \right)_s ds, \quad m \geq 0, \]  
\[ \mathcal{\tilde{U}}_{m+1}(x, y; t; r) = \mathcal{\tilde{U}}_m(x, y; s; r) + \int_0^t \lambda \left( \left( \mathcal{\tilde{U}}_m(x, y; s; r) \right)_s - \frac{\gamma^2}{2} \left( \mathcal{\tilde{U}}_m(x, y; s; r) \right)_{xx} \right)_s ds, \quad m \geq 0. \]

Making the above correction functional stationary, we have

\[ \delta \mathcal{U}_{m+1}(x, y; t; r) = \delta \mathcal{U}_m(x, y; t; r) + \delta \int_0^t \lambda \left( \left( \mathcal{U}_m(x, y; s; r) \right)_s - \frac{\gamma^2}{2} \left( \mathcal{U}_m(x, y; s; r) \right)_{xx} \right)_s ds, \quad m \geq 0, \]  
\[ \delta \mathcal{\tilde{U}}_{m+1}(x, y; t; r) = \delta \mathcal{\tilde{U}}_m(x, y; s; r) + \delta \int_0^t \lambda \left( \left( \mathcal{\tilde{U}}_m(x, y; s; r) \right)_s - \frac{\gamma^2}{2} \left( \mathcal{\tilde{U}}_m(x, y; s; r) \right)_{xx} \right)_s ds, \quad m \geq 0. \]

Similarly, the stationary conditions are given by

\[ 1 + \lambda |_{s=t}, \quad \lambda'(s) = 0, \]

and thus the Lagrange multiplier

\[ \lambda = -1. \]

Substituting the value of the Lagrange multiplier into equations (4.35) and (4.36) gives the iteration formulas

\[ \mathcal{U}_{m+1}(x, y; t; r) = \mathcal{U}_m(x, y; t; r) - \int_0^t \left[ \left( \mathcal{U}_m(x, y; s; r) \right)_s - \frac{\gamma^2}{2} \left( \mathcal{U}_m(x, y; s; r) \right)_{xx} \right] ds, \]
\[ \mathcal{\tilde{U}}_{m+1}(x, y; t; r) = \mathcal{\tilde{U}}_m(x, y; s; r) - \int_0^t \left[ \left( \mathcal{\tilde{U}}_m(x, y; s; r) \right)_s - \frac{\gamma^2}{2} \left( \mathcal{\tilde{U}}_m(x, y; s; r) \right)_{xx} \right] ds. \]

Using the iteration formulas of equations (4.39) and (4.40) with the initial conditions

\[ \hat{u}_0(x, y; t) = \hat{u}(x, y, t) = [(1 + 2t)^n, (5 - 2t)^n] \quad \text{or} \quad y^2, \quad n = 1, 2, 3, \ldots, \]

we get the following suc-
cessive approximations:

\[
\begin{aligned}
\phi_1(x, y, t; r) &= (1 + 2r)^n - (tx^2 + y^2), \\
\phi_2(x, y, t; r) &= (1 + 2r)^n - \left[ tx^2 + \left( 1 + \frac{t^2}{2!} \right) y^2 \right], \\
\phi_3(x, y, t; r) &= (1 + 2r)^n - \left[ \left( t + \frac{t^3}{3!} \right) x^2 + \left( 1 + \frac{t^2}{2!} \right) y^2 \right], \\
&\vdots \\
\phi_m(x, y, t; r) &= (1 + 2r)^n - \left[ \left( t + \frac{t^3}{3!} + \cdots + \frac{t^{2m+1}}{(2m+1)!} \right) x^2 \\
&\quad + \left( 1 + \frac{t^2}{2!} + \cdots + \frac{t^{2m}}{(2m)!} \right) y^2 \right],
\end{aligned}
\] (4.41)

and

\[
\begin{aligned}
\overline{\phi}_1(x, y, t; r) &= (5 - 2r)^n - (tx^2 + y^2), \\
\overline{\phi}_2(x, y, t; r) &= (5 - 2r)^n - \left[ tx^2 + \left( 1 + \frac{t^2}{2!} \right) y^2 \right], \\
\overline{\phi}_3(x, y, t; r) &= (5 - 2r)^n - \left[ \left( t + \frac{t^3}{3!} \right) x^2 + \left( 1 + \frac{t^2}{2!} \right) y^2 \right], \\
&\vdots \\
\overline{\phi}_m(x, y, t; r) &= (5 - 2r)^n - \left[ \left( t + \frac{t^3}{3!} + \cdots + \frac{t^{2m+1}}{(2m+1)!} \right) x^2 \\
&\quad + \left( 1 + \frac{t^2}{2!} + \cdots + \frac{t^{2m}}{(2m)!} \right) y^2 \right].
\end{aligned}
\] (4.42)

Consequently, the exact solution is given by

\[
\begin{aligned}
\underline{\phi}(x, y, t; r) &= (1 + 2r)^n - \left[ \lim_{m \to \infty} \left( t + \frac{t^3}{3!} + \cdots + \frac{t^{2m+1}}{(2m+1)!} \right) x^2 \\
&\quad + \left( 1 + \frac{t^2}{2!} + \cdots + \frac{t^{2m}}{(2m)!} \right) y^2 \right], \\
\overline{\phi}(x, y, t; r) &= (5 - 2r)^n - \left[ \lim_{m \to \infty} \left( t + \frac{t^3}{3!} + \cdots + \frac{t^{2m+1}}{(2m+1)!} \right) x^2 \\
&\quad + \left( 1 + \frac{t^2}{2!} + \cdots + \frac{t^{2m}}{(2m)!} \right) y^2 \right],
\end{aligned}
\]

and we can rewrite the solution as follows:

\[
\hat{\phi}(x, y, t; r) = \left[ (1 + 2r)^n, (5 - 2r)^n \right] \otimes_{gh} (x^2 \sinh t + y^2 \cosh t), \quad 0 \leq r < 1.
\]
Example 4.3 We consider the following three-dimensional inhomogeneous fuzzy heat-like equation with variable coefficients:

\[
\tilde{u}_t = \tilde{\tau}(x,y,z) \oplus \frac{1}{36} \left[ x^2 \odot \tilde{u}_{xx} \oplus y^2 \odot \tilde{u}_{yy} \oplus z^2 \odot \tilde{u}_{zz} \right], \quad 0 < x,y,z < 1, t > 0, \tag{4.43}
\]

with the boundary conditions

\[
\begin{align*}
\tilde{u}(0,y,z,t) &= 0, & \tilde{u}(1,y,z,t) &= y^4 z^4 (\exp(t) - 1), \\
\tilde{u}(x,0,z,t) &= 0, & \tilde{u}(x,1,z,t) &= x^4 z^4 (\exp(t) - 1), \\
\tilde{u}(x,y,0,t) &= 0, & \tilde{u}(x,y,1,t) &= x^4 y^4 (\exp(t) - 1),
\end{align*} \tag{4.44}
\]

and the initial condition

\[
\tilde{u}(x,y,z,0) = \tilde{0}, \tag{4.45}
\]

where

\[
\tilde{\tau}(x,y,z;r) = (-1,0,1)^n \odot (xyz)^4, \\
= \left[(r-1)^n, (1-r)^n \right] \odot (xyz)^4, \quad 0 \leq r \leq 1 \ (n = 1, 2, 3, \ldots), \tilde{0} \in E^1.
\]

The parametric form of (4.43) is

\[
\begin{align*}
\underline{u}_t &= (r-1)^n x^4 y^4 z^4 + \frac{1}{36} \left[ x^2 \underline{u}_{xx} + y^2 \underline{u}_{yy} + z^2 \underline{u}_{zz} \right], \quad 0 < x,y,z < 1, t > 0, \tag{4.46} \\
\overline{u}_t &= (1-r)^n x^4 y^4 z^4 + \frac{1}{36} \left[ x^2 \overline{u}_{xx} + y^2 \overline{u}_{yy} + z^2 \overline{u}_{zz} \right], \quad 0 < x,y,z < 1, t > 0, \tag{4.47}
\end{align*}
\]

for \( r \in [0,1] \), where \( \underline{u} \) stands for \( u(x,y,z,t;r) \), and similarly for \( \overline{u} \).

Case [A]. Fuzzy Adomian decomposition method

Working equations (4.46) and (4.47) are

\[
\begin{align*}
L_t \underline{u} &= (r-1)^n x^4 y^4 z^4 + \frac{1}{36} \left[ x^2 \underline{u}_{xx} + y^2 \underline{u}_{yy} + z^2 \underline{u}_{zz} \right](r), \tag{4.48} \\
L_t \overline{u} &= (1-r)^n x^4 y^4 z^4 + \frac{1}{36} \left[ x^2 \overline{u}_{xx} + y^2 \overline{u}_{yy} + z^2 \overline{u}_{zz} \right](r). \tag{4.49}
\end{align*}
\]

Applying \( L_t^{-1} \) to equations (4.48) and (4.49), we obtain

\[
\begin{align*}
\underline{u}(x,y,z,t;r) &= (r-1)^n x^4 y^4 z^4 t + \frac{1}{36} L_t^{-1} \left[ x^2 \underline{u}_{xx} + y^2 \underline{u}_{yy} + z^2 \underline{u}_{zz} \right](r), \tag{4.50} \\
\overline{u}(x,y,z,t;r) &= (1-r)^n x^4 y^4 z^4 t + \frac{1}{36} L_t^{-1} \left[ x^2 \overline{u}_{xx} + y^2 \overline{u}_{yy} + z^2 \overline{u}_{zz} \right](r). \tag{4.51}
\end{align*}
\]
Using the fuzzy ADM series (3.27) and (3.28) for the solution \( \tilde{u}(x, y, z; t) = [u(x, y, z; t), \overline{u}(x, y, z; t)] \) gives

\[
\sum_{n=0}^{\infty} u_m(x, y, z; t; r) = (r - 1)^n x^4 y^4 z^4 t + \frac{1}{36} \left[ x^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z; t; r) \right)_{xx} + y^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z; t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z; t; r) \right)_{zz} \right],
\]

(4.52)

\[
\sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) = (1 - r)^n x^4 y^4 z^4 t + \frac{1}{36} \left[ x^2 \left( \sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) \right)_{xx} + y^2 \left( \sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) \right)_{zz} \right].
\]

(4.53)

We use the recursive relation

\[
\begin{align*}
    u_0(x, y, z; t; r) &= (r - 1)^n x^4 y^4 z^4 t, \\
    u_m(x, y, z; t; r) &= \frac{1}{36} L^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z; t; r) \right)_{xx} + y^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z; t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z; t; r) \right)_{zz} \right], \\
    m &\geq 0,
\end{align*}
\]

and

\[
\begin{align*}
    \overline{u}_0(x, y, z; t; r) &= (1 - r)^n x^4 y^4 z^4 t, \\
    \overline{u}_m(x, y, z; t; r) &= \frac{1}{36} L^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) \right)_{xx} + y^2 \left( \sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} \overline{u}_m(x, y, z; t; r) \right)_{zz} \right], \\
    m &\geq 0.
\end{align*}
\]

The first few components of \( \tilde{u}(x, y, z; t) = [u(x, y, z; t), \overline{u}(x, y, z; t)] \) are

\[
\begin{align*}
    u_0(x, y, z; t; r) &= (r - 1)^n x^4 y^4 z^4 t, \\
    u_1(x, y, z; t; r) &= \frac{1}{36} L^{-1} \left[ x^2 (u_0)_{xx} + y^2 (u_0)_{yy} + z^2 (u_0)_{zz} \right] = (r - 1)^n x^4 y^4 z^4 \frac{t^2}{2}, \\
    u_2(x, y, z; t; r) &= \frac{1}{36} L^{-1} \left[ x^2 (u_1)_{xx} + y^2 (u_1)_{yy} + z^2 (u_1)_{zz} \right] = (r - 1)^n x^4 y^4 z^4 \frac{t^3}{3!}, \\
    u_3(x, y, z; t; r) &= \frac{1}{36} L^{-1} \left[ x^2 (u_2)_{xx} + y^2 (u_2)_{yy} + z^2 (u_2)_{zz} \right] = (r - 1)^n x^4 y^4 z^4 \frac{t^4}{4!}, \\
    \vdots \\
    u_m(x, y, z; t; r) &= \frac{1}{36} L^{-1} \left[ x^2 (u_{m-1})_{xx} + y^2 (u_{m-1})_{yy} + z^2 (u_{m-1})_{zz} \right] = (r - 1)^n x^4 y^4 z^4 \frac{t^m}{m!},
\end{align*}
\]

(4.54)
where

\[
\begin{aligned}
\tilde{u}_0(x, y, z, t; r) &= (1 - r)^a x^4 y^4 z^4 t, \\
\tilde{u}_1(x, y, z, t; r) &= \frac{1}{36} L^{-1}_x \left( x^2 (\tilde{u}_0)_{xx} + y^2 (\tilde{u}_0)_{yy} + z^2 (\tilde{u}_0)_{zz} \right) = (1 - r)^a x^4 y^4 z^4 \frac{t^2}{2!}, \\
\tilde{u}_2(x, y, z, t; r) &= \frac{1}{36} L^{-1}_x \left( x^2 (\tilde{u}_1)_{xx} + y^2 (\tilde{u}_1)_{yy} + z^2 (\tilde{u}_1)_{zz} \right) = (1 - r)^a x^4 y^4 z^4 \frac{t^3}{3!}, \\
\tilde{u}_3(x, y, z, t; r) &= \frac{1}{36} L^{-1}_x \left( x^2 (\tilde{u}_2)_{xx} + y^2 (\tilde{u}_2)_{yy} + z^2 (\tilde{u}_2)_{zz} \right) = (1 - r)^a x^4 y^4 z^4 \frac{t^4}{4!}, \\
&\vdots \\
\tilde{u}_m(x, y, z, t; r) &= \frac{1}{36} L^{-1}_x \left( x^2 (\tilde{u}_m)_{xx} + y^2 (\tilde{u}_m)_{yy} + z^2 (\tilde{u}_m)_{zz} \right) = (1 - r)^a x^4 y^4 z^4 \frac{t^m}{m!}.
\end{aligned}
\]  

The solution for the inhomogeneous fuzzy heat-like equations (4.46) and (4.47) in a series form is given by

\[
\begin{aligned}
\tilde{u}(x, y, z, t; r) &= (r - 1)^a x^4 y^4 z^4 \left[ t + \frac{t^2}{2!} + \frac{t^3}{3!} + \frac{t^4}{4!} + \cdots + \frac{t^m}{m!} \right], \\
\tilde{u}(x, y, z, t; r) &= (1 - r)^a x^4 y^4 z^4 \left[ t + \frac{t^2}{2!} + \frac{t^3}{3!} + \frac{t^4}{4!} + \cdots + \frac{t^m}{m!} \right].
\end{aligned}
\]

The exact solution of equations (4.46) and (4.47) when \( \alpha = 1 \) is

\[
\tilde{u}(x, y, z, t; r) = \left[ (r - 1)^a, (1 - r)^a \right] \odot (x^4 y^4 z^4 (\exp(t) - 1)), \quad 0 \leq r \leq 1.
\]

Case [B]. Fuzzy variational iteration method

The correction functionals for equations (4.46) and (4.47) are

\[
\begin{aligned}
\tilde{u}_{m+1}(x, y, z, t; r) &= \tilde{u}_m(x, y, z, t; r) \\
&\quad + \int_0^t \lambda \left[ \left( \tilde{u}_m(x, y, z, s; r) \right)_s - (r - 1)^a (xyz)^4 - F \right] ds, \\
\tilde{v}_{m+1}(x, y, z, t; r) &= \tilde{v}_m(x, y, z, t; r) \\
&\quad + \int_0^t \lambda \left[ \left( \tilde{v}_m(x, y, z, s; r) \right)_s - (r - 1)^a (xyz)^4 - F \right] ds,
\end{aligned}
\]

where \( m \geq 0, \)

\[
\begin{aligned}
F &= \frac{1}{36} \left( x^2 (\tilde{u}_m(x, y, z, s; r))_{xx} + y^2 (\tilde{u}_m(x, y, z, s; r))_{yy} + z^2 (\tilde{u}_m(x, y, z, s; r))_{zz} \right), \\
\mathcal{F} &= \frac{1}{36} \left( x^2 (\tilde{v}_m(x, y, z, s; r))_{xx} + y^2 (\tilde{v}_m(x, y, z, s; r))_{yy} + z^2 (\tilde{v}_m(x, y, z, s; r))_{zz} \right),
\end{aligned}
\]

and \( \tilde{u}_0(x, y, z, s; r) = 0 \) for \( 0 \in E^1. \)
Making the above correction functionals stationary, we have

\[
\delta u_{m+1}(x, y, z; t; r) = \delta u_m(x, y, z; t; r) \\
+ \delta \int_0^t \lambda \left( \left( u_m(x, y, z; s; r) \right)_s - (r-1)n(xy^4 - F) \right) ds, \quad (4.58)
\]

\[
\delta \bar{u}_{m+1}(x, y, z; t; r) = \delta \bar{u}_m(x, y, z; t; r) \\
+ \delta \int_0^t \lambda \left( \left( \bar{u}_m(x, y, z; s; r) \right)_s - (1-r)n(xy^4 - F) \right) ds. \quad (4.59)
\]

Therefore the stationary conditions are defined by

\[
1 + \lambda |_{s=t}, \quad \lambda'(s) = 0,
\]

and the Lagrange multiplier is

\[
\lambda = -1.
\]

Substituting the value of the Lagrange multiplier into equations (4.56) and (4.57), we obtain the iteration formulas

\[
\begin{align*}
\tilde{u}_{m+1}(x, y, z; t; r) &= \tilde{u}_m(x, y, t; r) \\
&\quad - \int_0^t \left( \left( \tilde{u}_m(x, y, z; s; r) \right)_s - (r-1)n(xy^4 - F) \right) ds, \quad (4.60)
\end{align*}
\]

\[
\begin{align*}
\bar{u}_{m+1}(x, y, z; t; r) &= \bar{u}_m(x, y, t; r) \\
&\quad - \int_0^t \left( \left( \bar{u}_m(x, y, z; s; r) \right)_s - (1-r)n(xy^4 - F) \right) ds. \quad (4.61)
\end{align*}
\]

Using the fuzzy iteration formulad (4.60) and (4.61) with initial approximation: \( \tilde{u}_0(x, y, t; r) = \tilde{u}(x, y, t; r) = \tilde{0} \), we have the following successive approximations:

\[
\begin{align*}
\tilde{u}_1(x, y, z; t; r) &= (r-1)ntx^4y^4z^4, \\
\tilde{u}_2(x, y, z; t; r) &= (r-1)n \left[ t + \frac{t^2}{2!} \right] x^4y^4z^4, \\
\tilde{u}_3(x, y, z; t; r) &= (r-1)n \left[ t + \frac{t^2}{2!} + \frac{t^3}{3!} \right] x^4y^4z^4, \\
\vdots \\
\tilde{u}_m(x, y, z; t; r) &= (r-1)n \left[ t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots + \frac{t^m}{m!} \right] x^4y^4z^4, \quad (4.62)
\end{align*}
\]
Thereafter the exact solution is given by

\[ u(x, y, z, t; r) = (r - 1)\lim_{m \to \infty} \left[ t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots + \frac{t^m}{m!}\right] x^4 y^4 z^4, \]

and we can rewrite the solution as follows:

\[ \tilde{u}(x, y, z, t; r) = \left[ (r - 1)^n, (1 - r)^n \right] \odot \left( x^4 y^4 z^4 \left( \exp(t - 1) \right) \right), \quad 0 \leq r \leq 1. \]

**Example 4.4** We consider the following one-dimensional fuzzy wave-like equation with variable coefficients:

\[ \ddot{u}(x, 0) = \left[ (r^n, (2 - r)^n) \right] \odot x, \quad \ddot{u}_t(x, 0) = \left[ (r^n, (2 - r)^n) \right] \odot x^2, \]

where \( n = 1, 2, 3, \ldots \).

The parametric form of (4.64) is

\[ \ddot{u}(x, 0) = \frac{1}{2} x^2 \dddot{u}, \quad 0 < x < 1, t > 0, \]

and similar for \( \dddot{u} \).

**Case [A]. Fuzzy Adomian decomposition method**

Applying the fuzzy ADM to equations (4.67) and (4.68) yields

\[ L_t \dddot{u} = \frac{1}{2} x^2 \dddot{u}, \quad 0 < x < 1, \]
\[ L_t u = \frac{1}{2} x^2 u_{xx}, \quad 0 < x < 1, \quad (4.70) \]

where \( L_t \) and \( L_t^{-1} \) are the second differential operator and its inverse, respectively. We have
the twofold integral operator
\[ L_t^{-1} (\cdot) = \int_0^t \int_0^t (\cdot) ds \, ds. \quad (4.71) \]

Using equations (4.69) and (4.70) and initial conditions (4.66), we get
\[ u(x, t; r) = r^n (x + x^2 t) + \frac{1}{2} L_t^{-1} (x^2 u_{xx}), \quad (4.72) \]
\[ \bar{u}(x, t; r) = (2 - r)^n (x + x^2 t) + \frac{1}{2} L_t^{-1} (x^2 \bar{u}_{xx}). \quad (4.73) \]

Using the fuzzy ADM series (3.40) and (3.41) for \( \tilde{u}(x, t; r) = [u(x, t; r), \bar{u}(x, t; r)] \) in (4.72) and (4.73), we obtain
\[ \sum_{m=0}^{\infty} u_m(x, t; r) = r^n (x + x^2 t) + \frac{1}{2} L_t^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} u_m(x, t; r) \right)_{xx} \right], \quad (4.74) \]
\[ \sum_{m=0}^{\infty} \bar{u}_m(x, t; r) = (2 - r)^n (x + x^2 t) + \frac{1}{2} L_t^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} \bar{u}_m(x, t; r) \right)_{xx} \right]. \quad (4.75) \]

Next, use the recursive relations
\[ u_0(x, t; r) = r^n (x + x^2 t), \]
\[ u_{m+1}(x, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 (u_m)_{xx} \right], \quad m \geq 0, \quad (4.76) \]
and
\[ \bar{u}_0(x, t; r) = (2 - r)^n (x + x^2 t), \]
\[ \bar{u}_{m+1}(x, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 (\bar{u}_m)_{xx} \right], \quad m \geq 0. \quad (4.77) \]

We have the first few components
\[ u_0(x, t; r) = r^n (x + x^2 t), \]
\[ u_1(x, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 (u_0)_{xx} \right] = r^n x^2 \frac{t^3}{3!}, \]
\[ u_2(x, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 (u_1)_{xx} \right] = r^n x^2 \frac{t^5}{5!}, \]
\[ u_3(x, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 (u_2)_{xx} \right] = r^n x^2 \frac{t^7}{7!}, \]
\[ \vdots \]
and

\[ \begin{align*}
\bar{u}_0(x, t; r) &= (2 - r)^n (x + x^2 t), \\
\bar{u}_1(x, t; r) &= \frac{1}{2} L^{-1} [x^2 (\bar{u}_0)_{xx}] = (2 - r)^n x^2 t^3 \\
\bar{u}_2(x, t; r) &= \frac{1}{2} L^{-1} [x^2 (\bar{u}_1)_{xx}] = (2 - r)^n x^2 t^5 \\
\bar{u}_3(x, t; r) &= \frac{1}{2} L^{-1} [x^2 (\bar{u}_2)_{xx}] = (2 - r)^n x^2 t^7 \\
\vdots &
\end{align*} \]

The solution in a fuzzy ADM series form is given by

\[ u(x, t; r) = r^n \left[ x + x^2 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} + \cdots \right) \right], \]

\[ \bar{u}(x, t; r) = (2 - r)^n \left[ x + x^2 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} + \cdots \right) \right], \]

and thus the exact solution is

\[ \tilde{u}(x, t; r) = \left[ r^n, (2 - r)^n \right] \odot (x + x^2 \sinh t), \quad 0 \leq r \leq 1. \]

Case [B]. Fuzzy variational iteration method

The correction functionals for equations (4.67) and (4.68) are

\[ \begin{align*}
\delta \bar{u}_{m+1}(x, t; r) &= \delta \bar{u}_m(x, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, s; r) \right)_{ss} - \frac{x^2}{2} \left( \bar{u}_m(x, s; r) \right)_{xx} \right] ds, \\
\delta \bar{u}_{m+1}(x, t; r) &= \delta \bar{u}_m(x, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, s; r) \right)_{ss} - \frac{x^2}{2} \left( \bar{u}_m(x, s; r) \right)_{xx} \right] ds,
\end{align*} \]

for \( m \geq 0 \), and \( \tilde{u}_0(x, t; r) = [r^n, (2 - r)^n] \odot (x + tx^2) \), where \( n = 1, 2, 3, \ldots \).

Making the above correction functional stationary, we have

\[ \begin{align*}
\delta \bar{u}_{m+1}(x, t; r) &= \delta \bar{u}_m(x, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, s; r) \right)_{ss} - \frac{x^2}{2} \left( \bar{u}_m(x, s; r) \right)_{xx} \right] ds, \\
\delta \bar{u}_{m+1}(x, t; r) &= \delta \bar{u}_m(x, t; r) + \int_0^t \lambda \left[ \left( \bar{u}_m(x, s; r) \right)_{ss} - \frac{x^2}{2} \left( \bar{u}_m(x, s; r) \right)_{xx} \right] ds.
\end{align*} \]

Therefore the stationary conditions are

\[ \lambda \big|_{s=t} = 0, \quad 1 - \lambda \big|_{s=t} = 0, \quad \lambda''(s) = 0. \]

By integrating the differential equation we have

\[ \lambda''(s) = 0 \quad \Rightarrow \quad \lambda(s) = c_1 s + c_2. \]
Using the boundary conditions, we arrive at
\[
1 - \lambda |_{s=t} = 0 \quad \Rightarrow \quad c_1 = 1, \\
\lambda |_{s=t} = 0 \quad \Rightarrow \quad c_2 = -t.
\]
Therefore the Lagrange multiplier is
\[
\lambda = s - t.
\]
Substituting the value of the Lagrange multiplier into functionals for equations (4.78) and (4.79) gives the iterative formulas
\[
\begin{align*}
u_{m+1}(x, t; r) &= \nu_m(x, t; r) + \int_0^t (s - t) \left[ \left( \nu_m(x, s; r) \right)_{ss} - \frac{x^2}{2} \left( \nu_m(x, s; r) \right)_{xx} \right] ds, \quad (4.82) \\
\bar{\nu}_{m+1}(x, t; r) &= \bar{\nu}_m(x, t; r) + \int_0^t (s - t) \left[ \left( \bar{\nu}_m(x, s; r) \right)_{ss} - \frac{x^2}{2} \left( \bar{\nu}_m(x, s; r) \right)_{xx} \right] ds. \quad (4.83)
\end{align*}
\]
Using the iterative formula of equations (4.82) and (4.83) with the initial approximation:
\[
\tilde{\nu}_0(x, t; r) = [r^n, (2 - r)^n] \odot (x + tx^2),
\]
we get the following successive approximations:
\[
\begin{align*}
u_1(x, t; r) &= r^n \left[ x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} \right) x^2 \right], \\
\nu_2(x, t; r) &= r^n \left[ x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} \right) x^2 \right], \\
\nu_3(x, t; r) &= r^n \left[ x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} + \cdots \right) x^2 \right], \\
&\vdots
\end{align*}
\]
and
\[
\begin{align*}ar{\nu}_1(x, t; r) &= (2 - r)^n \left[ x + \left( t + \frac{t^3}{3!} \right) x^2 \right], \\
\bar{\nu}_2(x, t; r) &= (2 - r)^n \left[ x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} \right) x^2 \right], \\
\bar{\nu}_3(x, t; r) &= (2 - r)^n \left[ x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} \right) x^2 \right], \\
&\vdots
\end{align*}
\]
So the exact solution is given in the closed form as follows:
\[
\begin{align*}
u(x, t; r) &= r^n \left[ \lim_{n \to \infty} x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} + \cdots \right) \right], \\
\bar{\nu}(x, t; r) &= (2 - r)^n \left[ \lim_{n \to \infty} x + \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} + \cdots \right) \right],
\end{align*}
\]
and we can rewrite the solution as follows:

\[ \tilde{u}(x, t; r) = [r^n, (2 - r)^n] \odot (x + x^2 \sinh t), \quad 0 \leq r \leq 1. \]

**Example 4.5** We consider the following two-dimensional fuzzy wave-like equation with variable coefficients:

\[ \tilde{u}_{tt} = \frac{1}{12} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}), \quad 0 < x, y < 1, t > 0, \quad (4.84) \]

with the Neumann boundary conditions

\[ \tilde{u}_x(0, y, t) = 0, \quad \tilde{u}_x(1, y, t) = 4 \cosh t, \]
\[ \tilde{u}_y(x, 0, t) = 0, \quad \tilde{u}_y(x, 1, t) = 4 \sinh t, \quad (4.85) \]

and the initial conditions

\[ \tilde{u}(x, y, 0) = [(0.2 + 0.2r)^n, (0.6 - 0.2r)^n] \odot x^4, \]
\[ \tilde{u}_t(x, y, 0) = [(0.2 + 0.2r)^n, (0.6 - 0.2r)^n] \odot y^4, \quad (4.86) \]

where \( n = 1, 2, 3, \ldots \)

The parametric form of (4.84) is

\[ \tilde{u}_{tt} = \frac{1}{12} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}), \quad 0 < x, y < 1, t > 0, \quad (4.87) \]
\[ \tilde{u}_{tt} = \frac{1}{12} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}), \quad 0 < x, y < 1, t > 0, \quad (4.88) \]

for \( r \in [0, 1] \), where \( \tilde{u} \) stands for \( \tilde{u}(x, y, t; r) \), and similarly for \( \tilde{u} \).

**Case [A]. Fuzzy Adomian decomposition method**

In an operator form, equations (4.87) and (4.88) are

\[ L_t \tilde{u} = \frac{1}{12} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}), \quad (4.89) \]
\[ L_t \tilde{u} = \frac{1}{12} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}). \quad (4.90) \]

Applying the inverse operator \( L^{-1} \) to both sides of equations (4.89) and (4.90), we obtain

\[ \tilde{u}(x, y, t; r) = (0.2 + 0.2r)^n (x^4 + y^4) + \frac{1}{12} L_t^{-1} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}), \quad (4.91) \]
\[ \tilde{u}(x, y, t; r) = (0.6 - 0.2r)^n (x^4 + y^4) + \frac{1}{12} L_t^{-1} (x^2 \tilde{u}_{xx} + y^2 \tilde{u}_{yy}). \quad (4.92) \]

Using the fuzzy ADM series for \( \tilde{u}(x, y, t; r) = [\tilde{u}(x, y, t; r), \tilde{u}(x, y, t; r)] \) gives

\[ \sum_{m=0}^{\infty} \tilde{u}_m(x, y, t; r) = (0.2 + 0.2r)^n (x^4 + y^4) + \frac{1}{12} \left[ x^2 \left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, t; r) \right) \right]_{xx} + y^2 \left( \sum_{m=0}^{\infty} \tilde{u}_m(x, y, t; r) \right)_{yy}. \]
\[\sum_{m=0}^{\infty} \mu_m(x, y; t; r) = (0.6 - 0.2r)^n \left(x^4 + y^4t\right) + \frac{1}{12} \int_1 \left[ x^2 \left( \sum_{m=0}^{\infty} \mu_m(x, y; t; r) \right)_{xx} + y^2 \left( \sum_{m=0}^{\infty} \mu_m(x, y; t; r) \right)_{yy} \right]. \]

After that, we use the recursive relations

\[\begin{align*}
\mu_0(x, y; t; r) &= (0.2 + 0.2r)^n \left(x^4 + y^4t\right), \\
\mu_{m+1}(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \mu_m \right)_{xx} + y^2 \left( \mu_m \right)_{yy} \right], \quad m \geq 0,
\end{align*}\]

and

\[\begin{align*}
\bar{\mu}_0(x, y; t; r) &= (0.6 - 0.2r)^n \left(x^4 + y^4t\right), \\
\bar{\mu}_{m+1}(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \bar{\mu}_m \right)_{xx} + y^2 \left( \bar{\mu}_m \right)_{yy} \right], \quad m \geq 0.
\end{align*}\]

The first few components of \(\tilde{u}(x, y; t; r)\) are given by

\[\begin{align*}
\tilde{u}_0(x, y; t; r) &= (0.2 + 0.2r)^n \left(x^4 + y^4t\right), \\
\tilde{u}_1(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \mu_0 \right)_{xx} + y^2 \left( \mu_0 \right)_{yy} \right] = (0.2 + 0.2r)^n \left(x^4 \frac{t^2}{2!} + y^4 \frac{t^1}{3!}\right), \\
\tilde{u}_2(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \mu_1 \right)_{xx} + y^2 \left( \mu_1 \right)_{yy} \right] = (0.2 + 0.2r)^n \left(x^4 \frac{t^4}{4!} + y^4 \frac{t^2}{5!}\right), \\
\tilde{u}_3(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \mu_2 \right)_{xx} + y^2 \left( \mu_2 \right)_{yy} \right] = (0.2 + 0.2r)^n \left(x^4 \frac{t^6}{6!} + y^4 \frac{t^2}{7!}\right), \\
\tilde{u}_4(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \mu_3 \right)_{xx} + y^2 \left( \mu_3 \right)_{yy} \right] = (0.2 + 0.2r)^n \left(x^4 \frac{t^8}{8!} + y^4 \frac{t^2}{9!}\right), \\
\vdots
\end{align*}\]

and

\[\begin{align*}
\bar{\mu}_0(x, y; t; r) &= (0.6 - 0.2r)^n \left(x^4 + y^4t\right), \\
\bar{\mu}_1(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \bar{\mu}_0 \right)_{xx} + y^2 \left( \bar{\mu}_0 \right)_{yy} \right] = (0.6 - 0.2r)^n \left(x^4 \frac{t^2}{2!} + y^4 \frac{t^3}{3!}\right), \\
\bar{\mu}_2(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \bar{\mu}_1 \right)_{xx} + y^2 \left( \bar{\mu}_1 \right)_{yy} \right] = (0.6 - 0.2r)^n \left(x^4 \frac{t^4}{4!} + y^4 \frac{t^3}{5!}\right), \\
\bar{\mu}_3(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \bar{\mu}_2 \right)_{xx} + y^2 \left( \bar{\mu}_2 \right)_{yy} \right] = (0.6 - 0.2r)^n \left(x^4 \frac{t^6}{6!} + y^4 \frac{t^3}{7!}\right), \\
\bar{\mu}_4(x, y; t; r) &= \frac{1}{12} \int_1 \left[ x^2 \left( \bar{\mu}_3 \right)_{xx} + y^2 \left( \bar{\mu}_3 \right)_{yy} \right] = (0.6 - 0.2r)^n \left(x^4 \frac{t^8}{8!} + y^4 \frac{t^3}{9!}\right), \\
\vdots
\end{align*}\]
Therefore the solution in a series form is

\[
\begin{align*}
    u(x, y, t; r) &= (0.2 + 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots \right) \right], \\
    u(x, y, t; r) &= (0.6 - 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots \right) \right],
\end{align*}
\]

and thus the exact solution is

\[
\tilde{u}(x, y, t; r) = \left[ (0.2 + 0.2r)^n, (0.6 - 0.2r)^n \right] \odot (x^4 \cosh t + y^4 \sinh t), \quad 0 \leq r \leq 1.
\]

Case [B]. Fuzzy variational iteration method

The correction functionals for equations (4.87) and (4.88) with the initial conditions (4.86) are

\[
\begin{align*}
    u_{m+1}(x, y, t; r) &= u_m(x, y, t; r) + \int_0^t \lambda \left[ \left( u_m(x, y, s; r) \right)_s - \frac{x^2}{12} \left( u_m(x, y, s; r) \right)_{xx} ight] ds, \\
    \n \end{align*}
\]

\[
\begin{align*}
    u_{m+1}(x, y, t; r) &= u_m(x, y, t; r) + \int_0^t \lambda \left[ \left( u_m(x, y, s; r) \right)_s - \frac{x^2}{12} \left( u_m(x, y, s; r) \right)_{xx} \right] ds, \\
    \n \end{align*}
\]

Making the above correction functional stationary, we have

\[
\begin{align*}
    \delta u_{m+1}(x, y, t; r) &= \delta u_m(x, y, t; r) + \delta \int_0^t \lambda \left[ \left( u_m(x, y, s; r) \right)_s - \frac{x^2}{12} \left( u_m(x, y, s; r) \right)_{xx} \right] ds, \\
    \n \end{align*}
\]

\[
\begin{align*}
    \delta u_{m+1}(x, y, t; r) &= \delta u_m(x, y, t; r) + \delta \int_0^t \lambda \left[ \left( u_m(x, y, s; r) \right)_s - \frac{x^2}{12} \left( u_m(x, y, s; r) \right)_{xx} \right] ds, \\
    \n \end{align*}
\]

Similarly, the stationary conditions are obtained by

\[
\lambda |_{s=t} = 0, \quad 1 - \lambda |_{s=t} = 0, \quad \lambda ''(s) = 0,
\]

and we obtain the Lagrange multiplier

\[
\lambda = s - t.
\]
Substituting the value of the Lagrange multiplier into equations (4.93) and (4.94) gives the iteration formulas

\[
\begin{align*}
\tilde{u}_{m+1}(x, y, t; r) &= \tilde{u}_m(x, y, t; r) + \int_0^t (s - t) \left[ \frac{1}{12} \left( \tilde{u}_m(x, y, s; r) \right)_{xx} \right. \\
&\quad - \frac{y^2}{12} \left( \tilde{u}_m(x, y, s; r) \right)_{yy} \bigg] ds, \\
\tilde{\pi}_{m+1}(x, y, t; r) &= \tilde{\pi}_m(x, y, t; r) + \int_0^t (s - t) \left[ \frac{1}{12} \left( \tilde{\pi}_m(x, y, s; r) \right)_{xx} \right. \\
&\quad - \frac{y^2}{12} \left( \tilde{\pi}_m(x, y, s; r) \right)_{yy} \bigg] ds.
\end{align*}
\] (4.97)

Using the iteration formula of equations (4.97) and (4.98) with the initial approximation \(\tilde{u}_0(x, y, t; r) = \tilde{u}(x, y, t; r) = [(0.2 + 0.2r)^n, (0.6 - 0.2r)^n] \odot (x^4 + t^4)\), where \(n = 1, 2, 3, \ldots\), we get the following successive approximations:

\[
\begin{align*}
\tilde{u}_1(x, y, t; r) &= (0.2 + 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} \right) + y^4 \left( t + \frac{t^3}{3!} \right) \right], \\
\tilde{u}_2(x, y, t; r) &= (0.2 + 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} \right) \right], \\
\tilde{u}_3(x, y, t; r) &= (0.2 + 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \frac{t^6}{6!} \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} \right) \right], \\
&\vdots
\end{align*}
\] (4.99)

and

\[
\begin{align*}
\tilde{\pi}_1(x, y, t; r) &= (0.6 - 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} \right) + y^4 \left( t + \frac{t^3}{3!} \right) \right], \\
\tilde{\pi}_2(x, y, t; r) &= (0.6 - 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} \right) \right], \\
\tilde{\pi}_3(x, y, t; r) &= (0.6 - 0.2r)^n \left[ x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \frac{t^6}{6!} \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} \right) \right], \\
&\vdots
\end{align*}
\] (4.100)

Consequently, the exact solution is given by

\[
\begin{align*}
\tilde{u}(x, y, t; r) &= (0.2 + 0.2r)^n \left[ \lim_{m \to \infty} x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots \right) \right], \\
\tilde{\pi}(x, y, t; r) &= (0.6 - 0.2r)^n \left[ \lim_{m \to \infty} x^4 \left( 1 + \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots \right) + y^4 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots \right) \right],
\end{align*}
\]

and we can rewrite the solution as follows:

\[
\tilde{u}(x, y, t; r) = [(0.2 + 0.2r)^n, (0.6 - 0.2r)^n] \odot (x^4 \cosh t + y^4 \sinh t), \quad 0 \leq r \leq 1.
\]
Example 4.6 We consider the following three-dimensional fuzzy wave-like equation with variable coefficients:

\[ \ddot{u}_{tt} = (x^2 + y^2 + z^2) \oplus \frac{1}{2} \left[ x^2 \odot \ddot{u}_{xx} + y^2 \odot \ddot{u}_{yy} + z^2 \odot \ddot{u}_{zz} \right], \]

where \( 0 < x, y, z < 1, t > 0, \) (4.101)

with the boundary conditions

\[ \ddot{u}(0, y, z, t) = y^2 \left( e^t - 1 \right) + z^2 \left( e^{-t} - 1 \right), \]
\[ \ddot{u}(1, y, z, t) = (1 + y^2) \left( e^t - 1 \right) + z^2 \left( e^{-t} - 1 \right), \]
\[ \ddot{u}(x, 0, z, t) = x^2 \left( e^t - 1 \right) + z^2 \left( e^{-t} - 1 \right), \]
\[ \ddot{u}(x, 1, z, t) = (x^2 + y^2) \left( e^t - 1 \right) + (e^{-t} - 1), \]

and the initial condition

\[ \ddot{u}(x, y, z, 0) = \ddot{0}, \quad \ddot{u}_t(x, y, z, 0) = \left[ 0.5(0.5r) + \left( x^2 + y^2 - z^2 \right) \right] \]

where \( n = 1, 2, 3, \ldots, \) and \( \ddot{0} \in \mathbb{F}^3. \)

The parametric form of (4.101) is

\[ \ddot{u}_{tt} = (x^2 + y^2 + z^2) + \frac{1}{2} \left[ x^2 \dddot{u}_{xx} + y^2 \dddot{u}_{yy} + z^2 \dddot{u}_{zz} \right], \quad 0 < x, y, z < 1, t > 0, \] (4.104)

\[ \ddot{u}_{tt} = (x^2 + y^2 + z^2) + \frac{1}{2} \left[ x^2 \dddot{u}_{xx} + y^2 \dddot{u}_{yy} + z^2 \dddot{u}_{zz} \right], \quad 0 < x, y, z < 1, t > 0, \] (4.105)

for \( r \in [0, 1], \) where \( \ddot{u} \) stands for \( \ddot{u}(x, y, z; t; r), \) and similarly of \( \ddot{u}_t. \)

Case [A]. Fuzzy Adomian decomposition method

In an operator form, equations (4.104) and (4.105) are

\[ L_t \ddot{u} = \left[ \left( x^2 + y^2 + z^2 \right) \right] + \frac{1}{2} \left[ x^2 \ddot{u}_{xx} + y^2 \ddot{u}_{yy} + y^2 \ddot{u}_{zz} \right] \]

\[ L_t \ddot{u} = \left[ \left( x^2 + y^2 + z^2 \right) \right] + \frac{1}{2} \left[ x^2 \ddot{u}_{xx} + y^2 \ddot{u}_{yy} + y^2 \ddot{u}_{zz} \right], \] (4.106)

(4.107)

Applying \( L_t^{-1} \) to equations (4.106) and (4.107), we obtain

\[ \ddot{u}(x, y, z; t; r) = (0.5r)^n + \left[ \left( x^2 + y^2 \right) \left( t + \frac{t^2}{21} \right) + z^2 \left( -t + \frac{t^2}{21} \right) \right] \]

\[ + \frac{1}{2} L_t^{-1} \left[ x^2 \ddot{u}_{xx} + y^2 \ddot{u}_{yy} + y^2 \ddot{u}_{zz} \right] \]
\[ u(x, y, z, t; r) = (1 - 0.5r)^n + \left[ (x^2 + y^2) \left( -t + \frac{t^2}{2!} \right) \right] \\
+ \frac{1}{2} L_t^{-1} \left[ x^2 \Pi_{xx} + y^2 \Pi_{yy} + z^2 \Pi_{zz} \right](r). \]

Using the fuzzy ADM series (3.40) and (3.41) for \( \tilde{u}(x, y, z, t; r) = [u(x, y, z, t; r), \Pi(x, y, z, t; r)] \), we obtain

\[
\sum_{m=0}^{\infty} u_m(x, y, z, t; r) = (0.5r)^m + \left[ \left( x^2 + y^2 \right) \left( -t + \frac{t^2}{2!} \right) \right] \\
+ \frac{1}{2} L_t^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z, t; r) \right)_{xx} \right. \\
+ y^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z, t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} u_m(x, y, z, t; r) \right)_{zz} \right], \quad (4.108)
\]

\[
\sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) = (1 - 0.5r)^m + \left[ \left( x^2 + y^2 \right) \left( -t + \frac{t^2}{2!} \right) \right] \\
+ \frac{1}{2} L_t^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) \right)_{xx} \right. \\
+ y^2 \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z, t; r) \right)_{zz} \right]. \quad (4.109)
\]

Subsequently, we use the recursive relations

\[
u_0(x, y, z, t; r) = (0.5r)^n + \left[ \left( x^2 + y^2 \right) \left( -t + \frac{t^2}{2!} \right) \right],
\]

\[
u_{m+1}(x, y, z, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} \nu_m(x, y, z; t; r) \right) \right]_{xx} \\
+ y^2 \left( \sum_{m=0}^{\infty} \nu_m(x, y, z; t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} \nu_m(x, y, z; t; r) \right)_{zz}, \quad m \geq 0,
\]

and

\[
\Pi_0(x, y, z, t; r) = (1 - 0.5r)^n + \left[ \left( x^2 + y^2 \right) \left( -t + \frac{t^2}{2!} \right) \right]
\]

\[
\Pi_{m+1}(x, y, z, t; r) = \frac{1}{2} L_t^{-1} \left[ x^2 \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z; t; r) \right) \right]_{xx} \\
+ y^2 \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z; t; r) \right)_{yy} + z^2 \left( \sum_{m=0}^{\infty} \Pi_m(x, y, z; t; r) \right)_{zz}, \quad m \geq 0.
\]
The first few components of \( \tilde{u}(x, y, z; t; r) \) are

\[
\begin{align*}
\tilde{u}_0(x, y, z; t; r) &= (0.5r)^n + \left[ (x^2 + y^2) \left( t + \frac{t^2}{2!} \right) + z^2 \left( -t + \frac{t^2}{2!} \right) \right], \\
\tilde{u}_1(x, y, z; t; r) &= \frac{1}{2} L_t^{-1} \left[ x^2 (\tilde{u}_0)_{xx} + y^2 (\tilde{u}_0)_{yy} + z^2 (\tilde{u}_0)_{zz} \right] \\
&= (0.5r)^n + \left[ (x^2 + y^2) \left( \frac{t^3}{3!} + \frac{t^4}{4!} \right) + z^2 \left( -\frac{t^3}{3!} + \frac{t^4}{4!} \right) \right], \\
\tilde{u}_2(x, y, z; t; r) &= \frac{1}{2} L_t^{-1} \left[ x^2 (\tilde{u}_1)_{xx} + y^2 (\tilde{u}_1)_{yy} + z^2 (\tilde{u}_1)_{zz} \right] \\
&= (0.5r)^n + \left[ (x^2 + y^2) \left( \frac{t^5}{5!} + \frac{t^6}{6!} \right) + z^2 \left( -\frac{t^5}{5!} + \frac{t^6}{6!} \right) \right], \\
\vdots
\end{align*}
\]

and

\[
\begin{align*}
\tilde{u}_0(x, y, z; t; r) &= (1 - 0.5r)^n + \left[ (x^2 + y^2) \left( t + \frac{t^2}{2!} \right) + z^2 \left( -t + \frac{t^2}{2!} \right) \right], \\
\tilde{u}_1(x, y, z; t; r) &= \frac{1}{2} L_t^{-1} \left[ x^2 (\tilde{u}_0)_{xx} + y^2 (\tilde{u}_0)_{yy} + z^2 (\tilde{u}_0)_{zz} \right] \\
&= (1 - 0.5r)^n + \left[ (x^2 + y^2) \left( \frac{t^3}{3!} + \frac{t^4}{4!} \right) + z^2 \left( -\frac{t^3}{3!} + \frac{t^4}{4!} \right) \right], \\
\tilde{u}_2(x, y, z; t; r) &= \frac{1}{2} L_t^{-1} \left[ x^2 (\tilde{u}_1)_{xx} + y^2 (\tilde{u}_1)_{yy} + z^2 (\tilde{u}_1)_{zz} \right] \\
&= (1 - 0.5r)^n + \left[ (x^2 + y^2) \left( \frac{t^5}{5!} + \frac{t^6}{6!} \right) + z^2 \left( -\frac{t^5}{5!} + \frac{t^6}{6!} \right) \right], \\
\vdots
\end{align*}
\]

The solution for the fuzzy wave-like equations (4.104) and (4.105) in a series form is given by

\[
\begin{align*}
\tilde{u}(x, y, z; t; r) &= (0.5r)^n + \left[ (x^2 + y^2) \left( t + \frac{t^2}{2!} + \frac{t^3}{3!} + \frac{t^4}{4!} + \cdots \right) \\
& \quad + z^2 \left( -t + \frac{t^2}{2!} - \frac{t^3}{3!} + \frac{t^4}{4!} - \cdots \right) \right], \\
\tilde{u}(x, y, z; t; r) &= (1 - 0.5r)^n + \left[ (x^2 + y^2) \left( t + \frac{t^2}{2!} + \frac{t^3}{3!} + \frac{t^4}{4!} + \cdots \right) \\
& \quad + z^2 \left( -t + \frac{t^2}{2!} - \frac{t^3}{3!} + \frac{t^4}{4!} - \cdots \right) \right].
\end{align*}
\]

The exact solution of equations (4.104) and (4.105) when \( \alpha = 1 \) is

\[
\tilde{u}(x, y, z; t; r) = \left[ (0.5r)^n, (1 - 0.5r)^n \right] \oplus \left( -\left( x^2 + y^2 + z^2 \right) + \left( x^2 + y^2 \right) e^t + z^2 e^{-t} \right),
\]

\[0 \leq r \leq 1.\]
Case [B]. Fuzzy variational iteration method

The correction functionals for equations (4.104) and (4.105) are

\[
\begin{align*}
\tilde{u}_{m+1}(x, y, z, t; r) &= u_m(x, y, z, t; r) + \int_0^t \lambda \left[ (u_m(x, y, z, s; r))_s \right] ds \\
&\quad - (x^2 + y^2 + z^2) + (0.5r)^n - F \right] ds, \\
\tilde{u}_{m+1}(x, y, z, t; r) &= \tilde{u}_m(x, y, z, t; r) + \int_0^t \lambda \left[ (\tilde{u}_m(x, y, z, s; r))_s \right] ds \\
&\quad - (x^2 + y^2 + z^2) + (1 - 0.5r)^n - \tilde{F} \right] ds,
\end{align*}
\]

for \( m \geq 0 \) and \( \tilde{u}_0(x, y, z, t; r) = [(0.5r)^n, (1 - 0.5r)^n] \oplus (x^2 + y^2 - z^2)t \), where \( n = 1, 2, 3, \ldots \).

Then

\[
\begin{align*}
F &= \frac{1}{2} (x^2 (u_m(x, y, z, s; r))_{xx} + (y^2 (u_m(x, y, z, s; r))_{yy} + (z^2 (u_m(x, y, z, s; r))_{zz}) \\
\tilde{F} &= \frac{1}{2} (x^2 (\tilde{u}_m(x, y, z, s; r))_{xx} + (y^2 (\tilde{u}_m(x, y, z, s; r))_{yy} + (z^2 (\tilde{u}_m(x, y, z, s; r))_{zz}).
\end{align*}
\]

Making the above correction functionals stationary, we have

\[
\begin{align*}
\delta u_{m+1}(x, y, z, t; r) &= \delta u_m(x, y, z, t; r) + \delta \int_0^t \lambda \left[ (u_m(x, y, z, s; r))_s \right] ds \\
&\quad - (x^2 + y^2 + z^2) + (0.5r)^n - F \right] ds, \\
\delta \tilde{u}_{m+1}(x, y, z, t; r) &= \delta \tilde{u}_m(x, y, z, t; r) + \delta \int_0^t \lambda \left[ (\tilde{u}_m(x, y, z, s; r))_s \right] ds \\
&\quad - (x^2 + y^2 + z^2) + (1 - 0.5r)^n - \tilde{F} \right] ds.
\end{align*}
\]

Similarly, the stationary conditions are obtained as

\[
\begin{align*}
\lambda |_{s=t} &= 0, \quad 1 - \lambda |_{s=t} = 0, \quad \lambda''(s) = 0,
\end{align*}
\]

and therefore the Lagrange multiplier is

\[
\lambda = s - t.
\]

Substituting the value of the Lagrange multiplier into functionals (4.112) and (4.113), we obtain the iteration formulas

\[
\begin{align*}
\tilde{u}_{m+1}(x, y, z, t; r) &= \tilde{u}_m(x, y, z, t; r) + \int_0^t (s - t) \left[ (\tilde{u}_m(x, y, z, s; r))_s \right] ds \\
&\quad - (x^2 + y^2 + z^2) + (0.5r)^n - F \right] ds, \\
\tilde{u}_{m+1}(x, y, z, t; r) &= u_m(x, y, z, t; r) + \int_0^t (s - t) \left[ (\tilde{u}_m(x, y, z, s; r))_s \right] ds \\
&\quad - (x^2 + y^2 + z^2) + (1 - 0.5r)^n - \tilde{F} \right] ds.
\end{align*}
\]
Therefore the exact solution is the closed form given by

\[
\begin{align*}
\bar{u}(x, y, z; t; r) &= (0.5r)^n + \left[ \lim_{m \to \infty} \left( t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots \right) \right] \left( x^2 + y^2 \right) \\
&\quad + z^2 \left[ -t + \frac{t^2}{2!} - \frac{t^3}{3!} + \cdots \right], \quad 0 \leq r \leq 1,
\end{align*}
\]

\[
\bar{u}(x, y, z; t; r) = (1 - 0.5r)^n + \left[ \lim_{m \to \infty} \left( t + \frac{t^2}{2!} + \frac{t^3}{3!} + \cdots \right) \right] \left( x^2 + y^2 \right) \\
&\quad + z^2 \left[ -t + \frac{t^2}{2!} - \frac{t^3}{3!} + \cdots \right],
\]

and we can rewrite the solution as follows:

\[
\tilde{u}(x, y, z; t; r) = \left[ (0.5r)^n, (1 - 0.5r)^n \right] \oplus \left( -\left( x^2 + y^2 + z^2 \right) + \left( x^2 + y^2 \right)e^t + z^2e^{-t} \right),
\]

Figures 1 and 2 are used to illustrate that the left-hand functions of the \(r\)-level set of \(\bar{u}\) (\(u\) lower) are always increasing functions of \(r\) and the right-hand functions of the \(r\)-level set of \(\tilde{u}\) (\(u\) upper) are always decreasing functions of \(r\) in the above examples.
5 Conclusions

In this paper, we investigated the fuzzy Adomian decomposition method (ADM) and fuzzy variational iteration method (VIM) and successfully applied to solving fuzzy heat-like and wave-like equations with variable coefficients under the $gH$-differentiability. To solve these equations, we first convert each fuzzy partial differential equation with its conditions into two crisp partial differential equations, and then we use ADM and VIM to numerically solve each crisp equation. We illustrate the proposed methods by some numerical examples and use figures to clarify the solutions. The results show that the fuzzy VIM reduces the volume of calculations, whereas tedious calculations are required in the fuzzy ADM; hence the iteration is direct and straightforward. Also, the results reveal that the methods are a very effective, convenient, and accurate mathematical tool for solving fuzzy heat-like and wave-like equations with variable coefficients.

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Authors’ contributions

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