Study of correlation of production and decay planes in $\pi \rightarrow 3\pi$ diffractive dissociation process on nuclei *

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Abstract

A large correlation of production and decay planes of ($\pi^-\pi^+\pi^-$)-system in dissociation of $\pi^-$-beam 40 GeV/c on nuclear targets was observed. The dependence of the correlation on atomic number, Feynmann variable and transversal momentum as well as on invariant mass of the pion triple and neutral pion pair were investigated. It was shown that the phenomenon has a clear dynamic origin and resembles the single spin asymmetry behavior.

1. Introduction

The measurement of polarization at high energies is a complicated problem since the observation of a spin or total angular momentum projection is a non-direct one as a rule and explore features such as angular distribution in secondary scattering or in decay process. For a strong interaction process parity and Lorentz invariance requires that at least three momenta of particles (either spinless or spin-averaged) in a final state were measured.

Some years ago the concept of handedness was introduced \(^1\) as a measure of polarization of parent partons (or decaying hadrons) \(^2\). It was defined as an asymmetry of a process probability $W$ with respect to a spatial component of an axial 4-vector $n_\mu \propto \epsilon_{\mu\nu\sigma\rho} k_1^\nu k_2^\sigma k_3^\rho$, where $k^\rho$ is 4-momentum of particle (or a system) in question ($k = k_1 + k_2 + k_3 + \cdots$), with respect to some direction $i$ ($n_i = ni$)

$$H_i = \frac{W(n_i > 0) - W(n_i < 0)}{W(n_i > 0) + W(n_i < 0)} = \alpha_i P_i,$$

which was shown to be proportional to polarization $P_i$ (at least for spin 1/2 and spin 1), provided the analyzing power $\alpha_i$ is not zero. The direction $i$ could be chosen as longitudinal ($L$) with respect to the momentum $k$ and as transversal ones ($T1$ or $T2$).

In the previous publication \(^3\) the attention was drawn to the fact that in diffractive production of pion triples \(^4\)

$$\pi^- + A \rightarrow (\pi^-\pi^+\pi^-) + A,$$

by $\pi^-$ beam 40 GeV/c from a nucleus $A$, a noticeable asymmetry with respect to the triple production plane (transversal handedness $H_{T1}$) was observed. This paper is devoted to further experimental investigation of this phenomenon. It includes a new information on the dependence of the transversal handedness on the variables:

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\(^1\) In fact, an idea similar to the handedness was earlier proposed in works \(^5\). Its application to certain heavy quark decays was studied in Ref. \(^6\). Similar technique was also studied in work \(^7\).
Atomic number of the target,  
Transversal momenta of the pion triple,  
Feynmann variable $X_F$ of the leading $\pi^-$,  
Invariant mass of the triple,  
Invariant mass of neutral pairs $\pi^+\pi^-$.  

Also the statistics was considerably increased.  

2. Definitions and notation  

For reaction (2), let us define the normal to the plane of production of a secondary pion triple $(\pi^- f \pi^+ s)$  
$$  N = (v_{3\pi} \times v_b) \quad (3)  $$

where $v_b = k_b/\epsilon_b$ and $v_{3\pi} = k_{3\pi}/\epsilon_{3\pi}$ are velocities of the initial $\pi^-$ beam and the center of mass of the triple in Lab. r.f. and indices $f$ and $s$ label fast and slow $\pi^-$'s. The normal to the "decay plane" of the triple in its center of mass is defined as  
$$  n = (v_f^- - v^+) \times (v_s^- - v^+) \quad (4)  $$

where $v_f^- (s)$ or $v^+$ are velocities of the fast (slow) $\pi^-$ or $\pi^+$.  

The transversal handedness according to (1) is  
$$  H_{T1} = \frac{W(Nn > 0) - W(Nn < 0)}{W(Nn > 0) + W(Nn < 0)}. \quad (5)  $$

Two other components of the handedness connected with $n \cdot v_{3\pi}$ and $n \cdot (v_{3\pi} \times N)$ are forbidden by the parity conservation in the strong interaction.  

3. Experimental results and discussion  

In this work the experimental material of Bologna–Dubna–Milan Collaboration for diffraction production of 40 GeV/c $\pi^-$ into three pions was used. The details of the experiment were presented in the works [6]. Notice here that the admixture of non-diffractive events in the used set of experimental data was less than 1%.  

The transversal handedness (3) was measured for a wide sample of nuclear targets: $Be$, $^{12}C$, $^{28}Si$, $^{48}Ti$, $^{63}Cu$, $^{107}Ag$, $^{181}Ta$ and $^{207}Pb$. The total number of selected events of pion triples with leading $\pi^-$ was about 250,000.  

The dependence of $H_{T1}$ on the atomic number $A$ is presented in Fig.1. One can see that the handedness systematically decreases with increasing $A$, which resembles a depolarization effect in multiple scattering. An argument in this favor is the decrease of the effect as, approximately, inverse nuclei radius.  

The value of the asymmetry (3), averaged over all nuclei is  
$$  H_{T1} = (5.96 \pm 0.21)\% \quad (6)  $$

Statistically, this is highly reliable verification of the existence of correlation of the triple production and decay planes in process (3).  

2It is easy to show that this quantity is in fact Lorentz-invariant.
The values of two other asymmetries with respect to correlations \( \mathbf{n} \cdot \mathbf{v}_{3\pi} \) and \( \mathbf{n} \cdot (\mathbf{v}_{3\pi} \times \mathbf{N}) \) was found to be comparable to zero from the same statistical material: \( H_L = (0.25 \pm 0.21)\% \) and \( H_{T2} = (0.43 \pm 0.21)\% \) respectively. This is by no means surprising, since they are forbidden by the parity conservation in process (2). Also they show the order of magnitude of systematic errors.

A natural question is to what extent the effect observed is due to the kinematics or apparatus influence, in particular, due to acceptance of the experimental setup where the events have been registered. For this aim the Monte-Carlo events of reaction (2) were generated with a constant mass spectrum of the \( 3\pi \)–system in the interval 0.6–2.5 GeV/c\(^2\) and with the exponential decrease of the cross section in \( t' = t - t_{\text{min}} \) with the slope (for beryllium) 40 (GeV/c\(^{-2}\)) found experimentally. This events were traced through the apparatus simulation with the same trigger conditions as in [3] and the same selection of events and show no transversal handedness \( H_{T1}^{\text{MC}} = (0.20 \pm 0.28)\% \)

\begin{equation}
H_{T1}^{\text{MC}} = (0.20 \pm 0.28)\%
\end{equation}

For two other asymmetries, forbidden by the parity conservation, the result was \((0.00 \pm 0.28)\%\) and \((-0.14 \pm 0.28)\%\), respectively. Thus, the effect (3) cannot be explained by the kinematics or apparatus influence.

To understand the nature of the effect observed, the dependence of the handedness (3) on the Feynmann variable \( X_F = k_f/k_b \) of the leading \( \pi^- \), on the invariant mass of the triple \( m_{3\pi} \) and its neutral subsystem \( m_{\pi^+\pi^-} \) and on the triple transversal momentum \( k_T \) was studied. From Fig.2 one can see that the handedness (3) increases with \( X_F \), which resembles the behavior of the single spin asymmetry (e.g. the pion asymmetry or the \( \Lambda \)-polarization [4]).

The dependence of \( H_{T1} \) on the triple invariant mass (Fig.3a) is especially interesting. It clearly indicates two different sources of \( H_{T1} \) with comparable contributions: a resonant and a non-resonant one. The resonance contribution is clearly seen at the mass of \( a_1(1260) \) and \( \pi_2(1670) \) region and by all means is due to a non-zero polarization of the resonances. The non-resonant background could also be polarized, provided that the \( 3\pi \) system is predominantly in a state with the total angular momentum \( J \neq 0 \), e.g. if a neutral pair \( m_{\pi^+\pi^-} \) was predominantly produced from \( \rho \)-decay. Some indication of this can be seen from Fig.3b. In this context, the growth of \( H_{T1} \) in the region of small \( m_{3\pi} \), i.e. in the region of small relative momenta of pions, looks quite intriguing.

A complicated picture of the \( k_T \)–dependence with a sharp deep at \( k_T = 0.05–0.07 \text{ GeV/c} \) (Fig.4) reflects by all means the fact of interference of the resonant and non-resonant processes in the triple production. With further increase of \( k_T \) the handedness increases which resembles the single spin asymmetry behavior too.

To check this assumption the events with invariant mass \( m_{3\pi} \) in the \( a_1 \) and \( \pi_2 \) resonance region \( 1.05–1.80 \text{ GeV} \) were excluded from further analysis. This however does not lead us to a definite conclusion since for \( Be \) and \( C \) the deep disappears but conserves for \( Si \) with some change of its form and width. The average value of the handedness stays at the same level \( 5–11\% \) with high statistical significance.

Notice also that in earlier study of reaction (2) at 4.5 GeV for the proton target at the hydrogen bubble chamber no angular dependence of the normal \( \mathbf{n} \) [4] was found just as in the Regge pole exchange model, which provides a reasonable description of that experiment [5].

In conclusion, a rather large handedness transversal to the production plane was definitely observed in the diffractive production of \( (\pi^-\pi^+\pi^-) \) triples in the \( \pi^- \)–beam dissociation region. The phenomenon has a clear dynamical origin and in some features resembles the single spin asymmetry behavior. For a more detailed study, a partial wave analysis of reaction (2) seems necessary for determination of different spin states contribution to the investigated effect.
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\[ \pi \rightarrow 3\pi \]
\[ \cdots \cdots \cdots \]
\[ (\pi^- \pi^+ \pi^-) \quad \pi^- \quad 40 \text{GeV/c} \]

\textbf{Figure captions}

Fig.1. The $A$-dependence of the handedness.

Fig.2. The handedness dependence on $X_F$ of the leading $\pi^-$.  

Fig.3. The handedness dependence on $m_{3\pi}$ (a) and $m_{\pi^+ \pi^-}$ (b).

Fig.4. The $k_T$-dependence of the handedness.
Fig. 1. The $A$-dependence of the handedness.

$$H = P_1 A^{P_2}$$

$$P_1 = 13.79 \pm 6.09$$

$$P_2 = -0.236 \pm 0.130$$

Fig. 2. The handedness dependence on $X_F$ of the leading $\pi^-$. 
Fig. 3. The handedness dependence on $m_{3\pi}$ (a) and $m_{\pi^+\pi^-}$ (b).
Fig. 4. The $k_T$-dependence of the handedness.