Interlaced processes on the circle

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Abstract. When two Markov operators commute, it suggests that we can couple two copies of one of the corresponding processes. We explicitly construct a number of couplings of this type for a commuting family of Markov processes on the set of conjugacy classes of the unitary group, using a dynamical rule inspired by the RSK algorithm. Our motivation for doing this is to develop a parallel programme, on the circle, to some recently discovered connections in random matrix theory between reflected and conditioned systems of particles on the line. One of the Markov chains we consider gives rise to a family of Gibbs measures on ‘bead configurations’ on the infinite cylinder. We show that these measures have determinantal structure and compute the corresponding space-time correlation kernel.

1. Introduction

When two Markov operators commute, it suggests that we can couple two copies of one of the corresponding processes. Such couplings have been described in \[10\] in a general context. In this paper, we explicitly construct a number of couplings of this type for a commuting family of Markov processes on the set of conjugacy classes of the unitary group, using a dynamical rule inspired by the Robinson-Schensted-Knuth (RSK) algorithm. Our motivation for doing this is to develop a parallel programme, on the circle, to some recently discovered connections in random matrix theory between reflected and conditioned systems of particles on the line (see, for example, \[2, 4, 19, 32, 40\]). The RSK algorithm is a combinatorial device which plays an important role in the representation theory of \(GL(n, \mathbb{C})\) and lies at the heart of these developments. We refer the reader to \[16\] for more background on

\[1991 \text{ Mathematics Subject Classification.} \ Primary \ 60J99, \ 60B15, \ 82B21; \ Secondary \ 05E10.\]
the combinatorics. One of the Markov chains we consider gives rise to a family of
Gibbs measures on ‘bead configurations’ on the infinite cylinder. This is related
to recent work [5, 6, 24] on planar and toroidal models. We will show that these
measures have determinantal structure and compute the corresponding space-time
correlation kernel.

We start with some motivation, and a flavour of some of the main results in
this paper. The following construction is closely related to the RSK algorithm [32, 30, 31]. Let \( E_1 = \mathbb{Z} \) and, for \( n \geq 2 \),
\[
E_n = \{ x \in \mathbb{Z}^n : x_1 < \cdots < x_n \}.
\]

The reader should think of an element \( x \in E_n \) as a configuration of \( n \) particles
on the integers located at positions \( x_1 < \cdots < x_n \). We say that a pair \((x,y) \in E_m \times E_{m+1}\) are interlaced, and write \( x \preceq y \), if \( y_j < x_j \leq y_{j+1} \) for all \( j \leq m \). A
(discrete) Gelfand-Tsetlin pattern of depth \( n \) is a collection \((x^1, x^2, \ldots, x^n)\) such
that \( x^m \in E_m \) for \( m \leq n \) and \( x^m \preceq x^{m+1} \) for \( 1 \leq m < n \).

Fix \( n \geq 2 \) and denote by \( GT_n \) the set of Gelfand-Tsetlin patterns of depth \( n \). Let \( w_1, w_2, \ldots \) be a sequence of independent random variables, each chosen according to
the uniform distribution on \( \{1, 2, \ldots, n\} \). Using these, we will construct a Markov
chain \((X(k), k \geq 0)\) with state space \( GT_n \), which evolves according to
\[
X(k+1) = g(X(k), w_{k+1}), \quad k \geq 0
\]
where \( g : GT_n \times \{1, \ldots, n\} \rightarrow GT_n \) is defined recursively as follows. Fix \( m < n \)
and let \((x,y) \in E_m \times E_{m+1}\) such that \( x \preceq y \). Let \( x' \in E_m \) such that, for some
\( j \leq m \), \( x'_i = x_i + \delta_{ij} \). The reader should have in mind \( m \) particles located at
positions \( x_1 < \cdots < x_m \), interlaced with another set of \( m+1 \) particles located
at positions \( y_1 < \cdots < y_{m+1} \). The first configuration \( x \) is updated by moving
the particle at position \( x_j \) one step to the right, that is, to position \( x_j + 1 \), giving
a new configuration \( x' \). This can be used to obtain an update \( y' \) to the second
configuration \( y \), obtained by moving the first available particle, strictly to the right
of position \( y_j \), which can be moved without breaking the interlacing constraint,
so that \( x' \preceq y' \). Such a particle is guaranteed to exist because the interlacing
constraint cannot be broken by the rightmost particle. In other words, the updated
configuration is given by \( y'_i = y_i + \delta_{ik} \) where \( k = \inf\{l > j : y_l + 1 < x_l\} \). Let
us write \( y' = \phi(x, y, x') \) where \( \phi \) is defined on an appropriate domain. Now, given \( m \leq n \) and a pattern \((x^1, \ldots, x^n) \in GT_n \) we define a new pattern

\[(y^1, \ldots, y^n) = g((x^1, \ldots, x^n), m)\]

as follows. First we set \( y^j = x^j \) for \( j < m \). Then we obtain \( y^m \) from the configuration \( x^m \) by moving the first available particle, starting from the particle at position \( x^m_1 \), which can be moved one step to the right without breaking the interlacing constraint. Finally, we define, for \( m \leq l < n \), \( y^{l+1} = \phi(x^l, x^{l+1}, y^l) \).

The Markov chain \( X \) has remarkable properties. For example, if we start from the initial pattern

\[(1) \quad X(0) = ((0), (-1, 0), (-2, -1, 0), \ldots, (-n + 1, \ldots, 0)), \]

then \((X^n(k), k \geq 0)\) is a Markov chain (with respect to its own filtration) with state space \( E_n \) and transition probabilities given by

\[(2) \quad P_n(x, y) = \begin{cases} \frac{1}{n} \frac{h(y)}{h(x)} & x \sim y \\ 0 & \text{otherwise} \end{cases} \]

where \( h \) is the Vandermonde function \( h(x) = \prod_{1 \leq j \leq n} (x_j - x_i) \) and the notation \( x \sim y \) means that the configuration \( y \) can be obtained from \( x \) by moving one particle one step to the right. The fact that \( P_n \) is a Markov transition kernel follows from the fact (see, for example, [25]) that \( h \) is positive on \( E_n \) and satisfies

\[\frac{1}{n} \sum_{x \sim y} h(y) = h(x).\]

More generally (see, for example, [30, 31]):

**Proposition 1.1.** Given \( x \in E_n \), if \( X(0) \) is chosen uniformly at random from the set of patterns \((x^1, \ldots, x^n) \) with \( x^n = x \), then \((X^n(k), k \geq 0)\) is a Markov chain started from \( x \) with transition probabilities given by \( P_n \). Moreover, for each \( T > 0 \), the conditional law of \( X(T) \), given \((X^n(k), T \geq k \geq 0)\), is uniformly distributed on the set of patterns \((x^1, \ldots, x^n) \) with \( x^n = X^n(T) \).

The connection with the RSK algorithm is the following: if we start from the initial pattern \( \mathbb{1} \), then the integer partition

\[(X^n_n(k), X^n_{n-1}(k) + 1, X^n_{n-2}(k) + 2, \ldots, X^n_1(k) + n - 1)\]
is precisely the shape of the tableau obtained when one applies the RSK algorithm, with column insertion, to the word \( w_1 w_2 \cdots w_k \). We refer to [30] for more details.

This construction clearly has a nested structure. If we consider the evolution of the last two rows \( X^{n-1} \) and \( X^n \) we see that we can construct a Markov chain with transition probabilities \( P_n \) from a Markov chain with transition probabilities \( P_{n-1} \) plus a ‘little bit of extra randomness’.

In the above construction we are thinking in terms of particles moving on a line. In random matrix theory, there are often strong parallels between natural measures on configurations of particles (or ‘eigenvalues’) on the line, and configurations of particles on the circle. It is therefore natural to ask if there is an analogue of the above construction for configurations of particles on the circle. The notion of interlacing carries over in the obvious way. However, in this setting, interlaced configurations should have the same number of particles, and the analogue of a Gelfand-Tsetlin pattern could potentially be an infinite object. Despite this fundamental difference between the two settings, there is indeed a natural analogue of the above ‘RSK dynamics’ and a natural analogue of Proposition 1.1. Consider the discrete circle with \( N \) positions which we label \( \{0, 1, \ldots, N-1\} \) in the anti-clockwise direction. The analogue of the Markov chain with transition matrix \( P_n \) is a Markov chain on the set \( C_n^N \) of configurations of \( n \) indistinguishable particles on the discrete circle with transition probabilities given by

\[
Q(x, y) = \begin{cases} 
\frac{x \Delta(y)}{n \Delta(x)} & x \nearrow y \\
0 & \text{otherwise}
\end{cases}
\]

where, similarly as before, the notation \( x \nearrow y \) means that the configuration \( y \) can be obtained from \( x \) by moving one particle one step anti-clockwise, and the function \( \Delta \) is again a Vandermonde function defined, for a configuration \( x \) which consists of \( n \) particles located at positions \( k_1, \ldots, k_n \), by

\[
\Delta(x) = \prod_{i<j} \left| e^{2\pi i k_j/N} - e^{2\pi i k_i/N} \right|.
\]

The constant \( c \) is chosen so that \( Q(x, \cdot) \) is a probability distribution; the fact that \( c \) can be chosen independently of \( x \) follows from the fact (see, for example, [25])
that $\Delta$ is a positive eigenvector of the matrix $1_{x \leq y}$, that is,

$$\sum_{x \leq y} \Delta(y) = \lambda \Delta(x),$$

for some $\lambda > 0$. A formula for the eigenvalue $\lambda$ can be found in [25].

In this setting we will say that a pair of configurations $(x, y)$ are interlaced, and write $x \preceq y$ as before, if there is a labelling of the particles such that $x$ consists of a particles located at positions $k_1, \ldots, k_n$, $y$ consists of a particles located at positions $l_1, \ldots, l_n$, $k_j < l_j \leq k_{j+1}$ for $j < n$ and either $k_n < l_n \leq k_1 + N$ or $k_n < l_n + N \leq k_1 + N$.

Let $(X(k), k \geq 0)$ be a Markov chain with state space $C_n^N$ and transition matrix $Q$. On the same probability space, without using any extra randomness, we can construct a second process $(Y(k), k \geq 0)$, also taking values in $C_n^N$, such that $X(k) \preceq Y(k)$ for all $k$. This is given as follows. For each $k > 0$, given $X(k)$, $Y(k)$ and $X(k+1)$ we obtain the configuration $Y(k+1)$ from $Y(k)$ by moving a particle one step anti-clockwise; this particle is chosen as follows. The transition from $X(k)$ to $X(k+1)$ involves one particle moving one step anti-clockwise; starting at the position of this particle, choose the first particle in the configuration $Y(k)$ that we come to, in an anti-clockwise direction, which can be moved by one step anti-clockwise without breaking the interlacing constraint.

The function $\Delta$ defined by (4) is also a positive (left and right) eigenvector of the matrix $1_{x \leq y}$. This follows, for example, from the discussion given at the end of Section 3. In particular,

$$\sum_{x \leq y} \Delta(x) = \gamma \Delta(y)$$

for some $\gamma > 0$ and we can define a Markov kernel on $C_n^N$ by

$$(5) \quad M(y, x) = \frac{1}{\gamma} \frac{\Delta(x)}{\Delta(y)} 1_{x \leq y}.$$ 

The analogue of Proposition 1.1 in this setting is the following:

**Proposition 1.2.** If $Y(0) = y$ and $X(0)$ is chosen at random according to the distribution $M(y, \cdot)$, then $(Y(k), k \geq 0)$ is a Markov chain started at $y$ with transition matrix $Q$. Moreover, for each $T > 0$, the conditional law of $X(T)$, given $(Y(k), T \geq k \geq 0)$, is given by $M(Y(T), \cdot)$. 
In the sequel we will present a number of variations of this result, firstly involving random walks with jumps in a continuous state space and secondly involving Brownian motion. Proposition 1.2 follows from results presented in Section 3 (see discussion towards the end of that section).

We will also study continuous analogues of the Markov chain with transition matrix $M$. These Markov chains also arise naturally in the context of a certain random walk on the unitary group which is obtained by taking products of certain random (complex) reflections, as studied for example in [15, 35]. This is described in Section 2 and taken as a starting point for the exposition that follows. The main point is that these Markov chains commute with each other and with the Dirichlet Laplacian on the set of conjugacy classes of the unitary group. In Sections 3 and 4, we present couplings which realise these commutation relations, first between interlaced random walks on the circle and later between interlaced Brownian motions. These couplings are precisely the variations of Proposition 1.2 mentioned above. Actually there are two natural couplings for the random walks and these correspond to dynamical rules inspired by the RSK algorithm with row, and column insertion, respectively. The couplings between interlaced Brownian motions can be thought of as a limiting case where the two types of coupling become equivalent. In Section 5, we consider a family of Markov chains which can be thought of as perturbations of a continuous analogue of the Markov chain with transition matrix given by [15]. These give rise to a natural family of Gibbs measures on ‘bead configurations’ on the infinite cylinder. This is a cylindrical analogue of the planar bead model studied in [5] and, in one special case (the ‘unperturbed’ case), can also be regarded as a cylindrical analogue of some natural measures on Gelfand-Tsetlin patterns related to ‘GUE minors’ [2, 23, 33, 14] (see also [7] for extensions to the other classical complex Lie algebras). We show that these measures have determinantal structure by first writing the restrictions of these measures to cylinder sets as products of determinants and then following the methodology of Johansson (see, for example, [22]) to compute the space-time correlation functions.

2. Markov processes on the conjugacy classes of the unitary group

Consider the group $U(n)$ of $n \times n$ unitary matrices, and denote by $C_n$ the set of conjugacy classes in $U(n)$. Each element of $C_n$ can be identified with an
unlabelled configuration of \( n \) points (eigenvalues) on the unit circle, which in turn can be identified with the Euclidean set \( D_n = \{ \theta \in \mathbb{R}^n : 0 \leq \theta_1 \leq \cdots \leq \theta_n < 2\pi \} \).

Denote by \( dx \) the image of Lebesgue measure under the latter identification, and by \( \mu \) the probability measure on \( C_n \) induced from Haar measure on \( U(n) \). Then

\[
\mu(dx) = (2\pi)^{-n} \Delta(x)^2 dx,
\]

where \( \Delta(x) \) is defined, for \( x = \{e^{i\theta_1}, \ldots, e^{i\theta_n}\} \), by

\[
(6) \quad \Delta(x) = \prod_{1 \leq l<m \leq n} |e^{i\theta_l} - e^{i\theta_m}|.
\]

The irreducible characters \( \chi_\lambda \) of \( U(n) \) are indexed by the set

\[
\Omega_n = \{ \lambda \in \mathbb{Z}^n : \lambda_1 \geq \lambda_2 \geq \cdots \geq \lambda_n \}.
\]

We assume that these are normalised, that is, \( \int |\chi_\lambda|^2 d\mu = 1 \) for all \( \lambda \in \Omega_n \).

Let \( x \in C_n \) and consider the random walk on \( U(n) \) which is constructed by multiplying together independent, randomly chosen elements from the conjugacy class \( x \). (By ‘a randomly chosen element from the conjugacy class \( x \)’ we mean a random element of the conjugacy class \( x \) which admits a representation of the form \( MD_x M^* \), where \( D_x \) is an arbitrary element of the conjugacy class \( x \) and \( M \) is a Haar-distributed, randomly chosen element of \( U(n) \).) The corresponding Markov kernel \( p_x(y, dz) \) can be interpreted as the law of \( D_x MD_y M^* \), where \( D_x \) and \( D_y \) are arbitrary elements of the conjugacy classes \( x \) and \( y \), respectively, and \( M \) is a Haar-distributed randomly chosen element of \( U(n) \). Write \( p_x f = \int p(\cdot, dz) f(z) \) for \( f \in L^2(C_n, \mu) \). Then, for each \( \lambda \in \Omega_n \),

\[
(7) \quad p_x \chi_\lambda = \frac{\chi_\lambda(x)}{d_\lambda} \chi_\lambda,
\]

where \( d_\lambda \) is the dimension of the representation corresponding to \( \lambda \) (see, for example, [13] Proposition 6.5.2)). Moreover, \( \mu \) is an invariant measure for the Markov kernel \( p_x \), that is,

\[
(8) \quad \int \mu(dy)p_x(y, \cdot) = \mu.
\]

The Markov kernels \( \{p_x, x \in C_n\} \) are the extreme points in the convex set \( \mathcal{M} \) of all Markov kernels on \( L^2(C_n, \mu) \) with the irreducible characters as eigenfunctions. (See, for example, [1]). All of the Markov kernels in \( \mathcal{M} \) have \( \mu \) as an invariant measure and, as operators on \( L^2(C_n, \mu) \), they commute. They correspond to random walks on \( U(n) \) such that the law of the increments is invariant under conjugation.
Such random walks on $U(n)$, and other classical compact groups, have been studied extensively in the literature (see, for example, [9, 8, 38, 34, 35]).

The case of interest in this paper is the random walk obtained by taking products of certain random (complex) reflections in $U(n)$ (see, for example, [15, 35]). More precisely, we take $x = \{ e^{ir}, 1, 1, \ldots, 1 \}$ in the above kernel, where $r \in (0, 2\pi)$. Let us write

$$p_r := p_x$$

for this case. A concrete description of this kernel can be given as follows (see [15] for details). For $a, b \in D_n$, write $a \preceq b$ if $a_1 \leq b_1 \leq \cdots \leq a_n \leq b_n$.

For $y = \{ e^{ia_1}, \ldots, e^{ian} \}$ and $z = \{ e^{ib_1}, \ldots, e^{ibn} \}$, where $a, b \in D_n$, write $y \preceq_r z$ if either $a \preceq b$ and $\sum_j (b_j - a_j) = r$, or $b \preceq a$ and $\sum_j (b_j - a_j) + 2\pi = r$. The measure $p_r(y, \cdot)$ is supported on the set

$$F_r(y) = \{ z \in C_n : y \preceq_r z \}.$$

This set can be identified with the disjoint union of a pair of $(n - 1)$-dimensional Euclidean sets

$$\{ b \in D_n : a \preceq b, \sum_j (b_j - a_j) = r \} \cup \{ b \in D_n : a \succeq b, \sum_j (b_j - a_j) = r - 2\pi \},$$

each of which can be endowed with $(n - 1)$-dimensional Lebesgue measure giving a natural measure on their union. The measure obtained on $F_r(y)$ via this identification can be extended to a measure $\nu_r(y, \cdot)$ on $C_n$ by setting $\nu_r(y, C_n \setminus F_r(y)) = 0$. The following identity can be deduced from [15 Lemma 2].

**Proposition 2.1.**

$$p_r(y, dz) = \frac{1}{\gamma_r} \frac{\Delta(z)}{\Delta(y)} \nu_r(y, dz),$$

where $\gamma_r = |1 - e^{ir}|^{(n-1)/(n-1)!}$.

In Section 5, we will present a determinantal formula for $\int_0^{2\pi} q^r \nu_r(y, dz) dr$, where $q > 0$ is a parameter, and use this to give an alternative proof of Proposition 2.1.
Another operator which will play a role in this paper is the Dirichlet Laplacian on the (closed) alcove

\[ A_n = \{ \theta \in \mathbb{R}^n : \theta_1 \leq \theta_2 \leq \cdots \leq \theta_n \leq \theta_1 + 2\pi \}. \]

Let \((Q_t)\) denote the transition semigroup of a standard Brownian motion conditioned never to exit \(A_n\). This is a Doob transform of the Brownian motion which is killed when it exits \(A_n\), via the positive eigenfunction

\[ h(\theta) = \prod_{1 \leq l < m \leq n} |e^{i\theta_l} - e^{i\theta_m}|, \]

We can identify \(C_n\) with \(\exp(iA_n)/\mathcal{C}_n\), where \(\mathcal{C}_n\) denotes the group of cyclic permutations which acts on \(A_n\) by permuting coordinates. It is known \([3]\) that the eigenfunctions of the induced semigroup \((\hat{Q}_t)\) on \(L^2(C_n, \mu)\) are given by the irreducible characters \(\{\chi_\lambda, \lambda \in \Omega_n\}\), which implies that \(\hat{Q}_t \in \mathcal{M}\), for each \(t > 0\). Note that the corresponding process on \(C_n\) can be thought of as \(n\) standard Brownian motions on the circle conditioned never to collide.

We will also consider discrete analogues of the above processes. Set

\[ A_n^N = (NA_n/2\pi) \cap \mathbb{Z}^n, \quad C_n^N = \exp(2\pi i A_n(N)/N)/\mathcal{C}_n, \quad \Omega_n^N = \{ \lambda \in \Omega_n : \lambda_1 \leq N-1 \}. \]

Much of the above discussion can be replicated in this setting, but for our purposes it suffices to make the following remark. Think of \(C_n^N\) as the set of configurations of \(n\) particles at distinct locations on the discrete circle with \(N\) positions. Consider the random walk in \(C_n^N\) where at each step a particle is chosen at random and moved one position anti-clockwise if that position is vacant; if it is not vacant the process is killed. It is known that the restriction of the function \(\Delta\) to \(C_n^N\) is the Perron-Frobenius eigenfunction for this sub-Markov chain. In fact, a complete set of eigenfunctions (with respect to the measure \(\Delta(x)^2\)) is given by the restrictions of the characters \(\{\chi_\lambda, \lambda \in \Omega_n^N\}\) to \(C_n^N\). (See, for example, \([25]\).) As we shall see later, the discrete analogues of the \(\nu_r\) (thought of as operators) commute with the transition kernel of this killed random walk and therefore share these eigenfunctions.

### 3. Couplings of interlaced random walks

The Markov chain on \(C_n\) with transition probabilities \(p_r\), defined by \([9]\), can be lifted to a Markov chain on \(A_n\), which is better suited to the constructions of this
To make this precise let us say \( x \) and \( x' \) belonging to \( A_n \) are \( r \)-interlaced, for some \( r \in (0, 2\pi) \), if

\[
x'_i \in [x_i, x_{i+1}] \text{ for } i = 1, 2, \ldots, n \text{ and } \sum_{i=1}^{n} (x'_i - x_i) = r,
\]

when we adopt the convention that \( x_{n+1} = x_1 + 2\pi \). In this case we will write \( x \preceq_r x' \). Define \( \pi : A_n \to C_n \) by \( \pi(x) = \{e^{ix_1}, \ldots, e^{ix_n}\} \). Denote by \( l_r(x, dx') \) the \((n-1)\)-dimensional Lebesgue measure on the set

\[
G_r(x) = \{x' \in A_n : x \preceq_r x'\}.
\]

Clearly the restriction of \( \pi \) to \( G_r(x) \) is injective, with \( \pi(G_r(x)) = F_r(\pi(x)) \). Moreover, for measurable \( B \subset A_n \), \( l_r(x, B) = \nu_r(\pi(x), \pi(B \cap G_r(x))) \). Thus, if we define, for measurable \( B \subset A_n \),

\[(14)\quad q_r(x, B) = p_r(\pi(x), \pi(B \cap G_r(x)))\]

then, by Proposition 2.1,

\[
q_r(x, B) = \int_{\pi(B \cap G_r(x))} p_r(\pi(x), dz)
= \int_{\pi(B \cap G_r(x))} \frac{1}{\gamma_r} \frac{\Delta(z)}{\Delta(\pi(x))} \nu_r(\pi(x), dz)
= \int_{B} \frac{1}{\gamma_r} \frac{\Delta(\pi(x'))}{\Delta(\pi(x))} l_r(x, dx'),
\]

and hence,

\[(15)\quad q_r(x, dx') = \frac{1}{\gamma_r} \frac{\Delta(\pi(x'))}{\Delta(\pi(x))} l_r(x, dx').\]

We will refer to a Markov chain with values in \( A_n \) and transition probabilities \( q_r \) as an \( r \)-interlacing random walk. As far as we are aware, such processes have not previously appeared in the literature. Note that since \( p_r p_s = p_s p_r \) we have \( q_r q_s = q_s q_r \) or, equivalently, \( l_r l_s = l_s l_r \).

The goal of this section is to construct, for given \( r, s \in (0, 2\pi) \), two different Markovian couplings \( (X(k), Y(k); k \geq 0) \) of a pair of \( r \)-interlacing random walks on \( A_n \), having the property that \( X(k) \) and \( Y(k) \) are \( s \)-interlaced for each \( n \) and moreover, for each \( l \geq 0 \), the trajectory \( (Y(k); 0 \leq k \leq l) \) will be a deterministic function of \( Y(0) \) together with the trajectory \( (X(k); 0 \leq k \leq l) \).
The existence of such couplings is suggested by the commutation relation \( q_r q_s = q_s q_r \), equivalently \( l_s l_s = l_r l_r \). For any \( u, v \in A_n \) consider the two sets \( \tau_{u,v} = \{ x \in A_n : u \preceq x \preceq v \} \) and \( \tau'_{u,v} = \{ y \in A_n : u \preceq y \preceq v \} \). If either is non-empty, then they both are, and in this case they are \((n-1)\)-dimensional polygons, and the relation \( l_s l_s = l_r l_r \) can be interpreted as saying these two polygons have the same \((n-1)\)-dimensional volume. In fact the two polygons are congruent. Define

\[
y = \phi_{u,v}(x)
\]

via

\[
y_i = \min(u_{i+1}, v_i) + \max(u_i, v_{i-1}) - x_i.
\]

It is easy to see that \( \phi_{u,v} \) is an isometry from \( \tau_{u,v} \) to \( \tau'_{u,v} \) using the facts that \( y_i \in [\max(u_i, v_{i-1}), \min(u_{i+1}, v_i)] \) if and only if \( x_i \in [\max(u_i, v_{i-1}), \min(u_{i+1}, v_i)] \), and

\[
\sum_{i=1}^{n} y_i = \sum_{i=1}^{n} (\min(u_{i+1}, v_i) + \max(u_i, v_{i-1}) - x_i) = \sum_{i=1}^{n} (u_i + v_{i-1} - x_i).
\]

**Proposition 3.1.** Let \((X(k); k \geq 0)\) be an \(r\)-interlacing random walk, starting from \(X(0)\) having the distribution \(q_r(y, dx)\) for some \(y \in A_n\), where \(q_r\) is defined by \((14)\). Let the process \((Y(k); k \geq 0)\) be given by \(Y(0) = y\) and

\[
Y(k + 1) = \phi_{Y(k),X(k+1)}(X(k)), \quad \text{for } k \geq 0,
\]

where \(\phi\) is defined by \((16)\). Then \((Y(k); k \geq 0)\) is distributed as an \(r\)-interlacing random walk starting from \(y\).

**Proof.** We prove by induction on \(m\) that the law of \((Y(1), \ldots, Y(m), X(m))\) is given by \(q_r(y, dy(1)) \cdots q_r(y(m-1), dy(m)) q_s(y(m), dx(m))\). Suppose this holds for some \(m\). Then, since \(Y(1), \ldots, Y(m)\) are measurable with respect to \(X(0), X(1), \ldots, X(m)\) the joint law of \((Y(1), \ldots, Y(m), X(m), X(m+1))\) is given by

\[
q_r(y, dy(1)) \cdots q_r(y(m-1), dy(m)) q_s(y(m), dx(m)) q_r(x(m), dx(m+1)).
\]

Equivalently we may say that the law of \((Y(1), \ldots, Y(m), X(m+1))\) is

\[
q_r(y, dy(1)) \cdots q_r(y(m-1), dy(m)) (q_s q_r)(y(m), dx(m+1))
\]
and that the conditional law of $X(m)$ given the same variables is uniform on $\tau Y(m) X(m+1)$. From the measure preserving properties of the maps $\phi_{u,v}$, it follows that, conditionally on $(Y(1), \ldots Y(m), X(m+1))$, $Y(m+1)$ is distributed uniformly on $\tau' Y(m) X(m+1)$, and the inductive hypothesis for $m+1$ follows from this and the commutation relation $q_s q_r = q_r q_s$. □

The dynamics of the coupled processes $(X(k), Y(k); k \geq 0)$ are illustrated in the following two diagrams, in which interlacing configurations $x = X(k)$ and $y = Y(k)$ are shown together with updated configurations $x' = X(k+1)$ and $y' = Y(k+1)$. It is natural to think of these as particle positions on (a portion of) the circle. For simplicity we consider an example where $x'$ and $x$ differ only in the $i$-th co-ordinate.

The configuration $y'$ is of course determined by $y$, $x$ and $x'$ together. The simplest possibility is shown above, in this case $i$th $y$ particle advances by the same amount as the $i$th $x$-particle. However should the $i$th $x$-particle advance beyond $y_{i+1}$ then it pushes the $(i + 1)$th $y$ particle along, whilst the increment in the position of the $i$th $y$ particle is limited to $y_{i+1} - x_i$.

The proof of Proposition 3.1 made use of only the measure preserving properties of the family of maps $\phi_{uv}$, and consequently we can replace it by another family of maps having the same property and obtain a different coupling of the same processes. The pushing interaction in the coupling constructed above is also seen in the dynamics induced on Gelfand-Tsetlin patterns by the RSK correspondence. There exists a variant of the RSK algorithm (with column insertion replacing the more common row insertion), in which pushing is replaced by blocking. We will
next describe a family $\psi_{uv}$ of measure-preserving maps that lead to a coupling with such a blocking interaction.

We recall first a version of the standard Skorohod lemma for periodic sequences.

**Lemma 3.2 (Skorohod).** Suppose that $(z_i; i \in \mathbb{Z})$ is $n$-periodic and satisfies $\sum_{i=1}^{n} z_i < 0$. Then there exists a unique pair of $n$-periodic sequences $(r_i; i \in \mathbb{Z})$ and $(l_i; i \in \mathbb{Z})$ such that

$$r_{i+1} = r_i + z_i + l_{i+1}$$

for all $i \in \mathbb{Z}$, with the additional properties $r_i \geq 0$, $l_i \geq 0$, and $l_i > 0 \implies r_i = 0$ for all $i \in \mathbb{Z}$.

A configuration $(x_1, x_2, \ldots, x_n)$ of $n$ points on the circle will be implicitly extended to a sequence $(x_i; i \in \mathbb{Z})$ satisfying $x_i + n = x_i + 2\pi$.

**Proposition 3.3.** Suppose that $u, v,$ and $x$ are three configurations on the circle with $u \preceq_s x \preceq_r v$, where $s > r$. Define an $n$-periodic sequence via

$$z_i = v_i - x_i - x_{i+1} + u_{i+1}$$

for $i \in \mathbb{Z}$, and let $(r, l)$ be the associated solution to the Skorohod problem. Set $y_i = x_i - l_i$, then $(y_1, \ldots, y_n)$ is configuration on the circle such that $u \preceq_r y \preceq_s v$.

**Proof.** Notice that $\sum_{i=1}^{n} z_i = r - s < 0$ so the appeal to the Skorohod construction is legitimate.

Summing the Skorohod equation gives $\sum_{i=1}^{n} z_i = -\sum_{i=1}^{n} l_i$. Consequently, using the definition of $y$,

$$\sum_{i=1}^{n} (y_i - u_i) = \sum_{i=1}^{n} (x_i - l_i - u_i) = s - \sum_{i=1}^{n} l_i = r.$$

Similarly $\sum_{i=1}^{n} (v_i - y_i) = s$.

To conclude we will verify the inequalities

$$\max(u_i, v_i - 1) \leq y_i \leq x_i \leq \min(u_{i+1}, v_i).$$

It is easily checked that $z_{i-1} \leq \min(x_i - u_i, x_i - v_i - 1)$. The solution of the Skorohod problem satisfies $0 \leq l_i \leq z_{i-1}$. Together with the definition of $y_i$, this gives the desired inequalities. \qed
By virtue of the preceding result we may define $\psi_{uv}: \tau_{uv} \rightarrow \tau'_{uv}$ via $\psi_{uv}(x) = y$.

**Proposition 3.4.** $\psi_{uv}$ is a measure-preserving bijection.

**Proof.** For $x \in A_n$, extended as before to a sequence $(x_i; i \in \mathbb{Z})$, let $x^\dagger$ be defined by $x^\dagger_i = -x_{-i}$ for $i \in \mathbb{Z}$. Observe that if $x \preceq_r y$, then $y^\dagger \preceq_r x^\dagger$. For $u, v, x$ and $y$ as in the previous proposition we will show that the application

$$(u, x, v) \mapsto (v^\dagger, y^\dagger, u^\dagger)$$

is an involution, and this implies in particular that $\psi_{uv}$ is invertible with $x^\dagger = \psi_{v^\dagger u^\dagger}(y^\dagger)$. To this end first note that $v^\dagger \preceq_r y^\dagger \preceq_r u^\dagger$, and hence it is meaningful to apply $\psi_{v^\dagger u^\dagger}$ to $y^\dagger$. We must prove that the resulting configuration $\tilde{x}$ say, is equal to $x^\dagger$.

We have $\tilde{x}_i = y^\dagger_i - \tilde{l}_i$, where $(\tilde{r}, \tilde{l})$ solves the Skorohod problem with data

$$\tilde{z}_i = u^\dagger_i - y^\dagger_i - y^\dagger_{i+1} + v^\dagger_i.$$ 

Since $\tilde{x}_i = y^\dagger_i - \tilde{l}_i = -x_{-i} + l_{-i} - \tilde{l}_i$, verifying that $\tilde{x} = x^\dagger$ boils down to checking that $\tilde{l}_i = l_{-i}$. For this it is suffices, by the uniqueness property of solutions to the Skorohod problem, to confirm that $r'_i = r_{-i}$ and $l'_i = l_{-i}$ solve the Skorohod problem with data $\tilde{z}$. Now $r'$ and $l'$ are non-negative and satisfy $l'_i > 0 \implies r'_i = 0$ because $r$ and $l$ have these properties. We just have to compute

$$r'_i + \tilde{z}_i + l'_{i+1} = r_{-i} + u^\dagger_i - y^\dagger_i - y^\dagger_{i+1} + v^\dagger_i + l_{-(i+1)}$$

$$= r_{-i} - u_{-i} + y_{-i} + y_{-(i+1)} - v_{-(i+1)} + l_{-(i+1)}$$

$$= r_{-i} - u_{-i} + x_{-i} - l_{-i} + x_{-(i+1)} - l_{-(i+1)} - v_{-(i+1)} + l_{-(i+1)}$$

$$= r_{-i} - z_{-(i+1)} - l_{-i}$$

$$= r_{-(i+1)} = r'_{i+1}.$$ 

This proves the involution property.

It remains to verify the measure-preserving property. The construction of $\psi_{uv}$ is such that it is evident that it is a piecewise linear mapping, and that its Jacobian is almost everywhere integer valued. Since the same applies to the inverse map constructed from $\psi_{v^\dagger u^\dagger}$ we conclude the Jacobian is almost everywhere $\pm 1$ valued. 

$\Box$
Finally the following proposition follows by exactly the same argument as Proposition 3.1.

**Proposition 3.5.** Suppose that $s > r$. Let $(X(k); k \geq 0)$ be an $r$-interlacing random walk, starting from $X(0)$ having the distribution $q_s(y, dx)$ for some given $y \in A_n$. Let the process $(Y(k); k \geq 0)$ be given by $Y(0) = y$ and

$$Y(k + 1) = \psi_{Y(k), X(k+1)}(X(k)),$$

for $k \geq 0$.

Then $(Y(k); k \geq 0)$ is distributed as an $r$-interlacing random walk starting from $y$.

Let us illustrate this coupling in a similar fashion to before.

In the simplest case shown above, the $(i + 1)$th $y$ particle advances by the same amount as the $i$th $x$-particle. However in the event that this would result in it passing the position $x_i$, then it is blocked at that point, and the unused part of the increment $(x'_i - x_i)$ is passed to the $(i + 2)$th $y$ particle as is shown beneath.

We can also consider discrete analogues of these constructions. Set $A_n^N = A_n \cap \mathbb{Z}^n/(2\pi N)$. Then, for $r \in \{1/(2\pi N), 2/(2\pi N), \ldots, (N - 1)/(2\pi N)\}$ we may define a Markovian transition kernel $q_r^N$ on $A_n^N$ as follows. Define

$$l_r^N(x, x') = \begin{cases} 1 & \text{if } z \leq_r x', \\ 0 & \text{otherwise}. \end{cases}$$

and let

$$h^N(x) = \prod_{1 \leq |i| < m \leq n} |e^{ix_i} - e^{ix_m}|. \quad (18)$$
where \( \bar{x}_j = \frac{N}{N+n}(x_j + j/(2\pi N)) \). Then, as we shall see below, \( l^N_r \), as a kernel on \( A^N_n \), admits \( h^N \) as a strictly positive eigenfunction, and consequently

\[
q^N_r(x, x') = \frac{1}{\gamma^N_r h^N(x)} l^N_r(x, x'),
\]

defines a Markovian transition kernel, where \( \gamma^N_r > 0 \) denotes the Perron-Frobenius eigenvalue of \( l^N_r \). We will call a Markov chain \( (X(k); k \geq 0) \) with these transition probabilities, an \( r \)-interlacing random walk on on \( A^N_n \).

In the special case \( r = 1/(2\pi N) \) the \( r \)-interlacing random walk on on \( A^N_n \) is closely related to non-colliding random walks on the circle, as considered in \( \text{[25]} \) for instance. Specifically if \( (X(k); k \geq 0) \) is the walk on \( A^N_n \) then the process given by

\[
(\psi_{i\bar{X}_1(k)}, \psi_{i\bar{X}_2(k)}, \ldots, \psi_{i\bar{X}_n(k)}),
\]

where \( \bar{X}_j(k) = \frac{N}{N+n}(X_j(k) + j/(2\pi N)), \) gives a process of \( n \) non-colliding random walks on the circle \( \{e^{ik/(2\pi(N+n)); k = 0, 1, 2, \ldots, N+n-1}\}. \)

Fix \( r \) and \( s \in \{1/(2\pi N), 2/(2\pi N), \ldots, (N-1)/(2\pi N)\} \). For \( u, v \in A^N_n \) we may consider the two sets \( \tau^N_{u,v} = \tau_{u,v} \cap \mathbb{Z}^n/(2\pi N) \) and \( \tau'_{u,v} = \tau'_{u,v} \cap \mathbb{Z}^n/(2\pi N) \), where \( \tau_{u,v} \) and \( \tau'_{u,v} \) are the polygons defined previously. The map \( \phi_{u,v} \) carries \( \mathbb{Z}^n/(2\pi N) \) into \( \mathbb{Z}^n/(2\pi N) \), and since we know it to be an isometry between \( \tau_{u,v} \) and \( \tau'_{u,v} \) it must therefore restrict to a bijection between \( \tau^N_{u,v} \) and \( \tau'^N_{u,v} \). In particular the cardinalities of these two sets are equal. Since \( l^N_{r,s}(u,v) = |\tau^N_{u,v}| = |\tau'^N_{u,v}| = l^N_{r,s}(u,v) \) we deduce that \( l^N_{r} \) and \( l^N_{s} \), and so \( q^N_{r} \) and \( q^N_{s} \) commute. This is useful in verifying our previous assertion that \( h^N \) is an eigenfunction of \( l^N_r \). The case \( r = 1 \) may be easily deduced by direct calculation, see \( \text{[25]} \) for a similar calculation. Moreover the statespace \( A^N_n \) is irreducible for \( l^N_r \) for \( r = 1 \) (though not in general), and thus the Perron-Frobenius eigenfunction of \( l^N_1 \) is unique up to multiplication by scalars. Consequently it follows from the commutation relation that \( h^N \) is an eigenfunction of \( l^N_r \) for every \( r \).

Using the fact that \( \phi_{u,v} \) is a bijection between \( \tau^N_{u,v} \) and \( \tau'^N_{u,v} \), we see that, substituting \( \tau^N_{u,v} \) for \( \tau_{u,v} \), \( q^N \) for \( q_r \) and so on, the proof of Proposition \( \text{[3.1]} \) applies verbatim in the case that \( X \) is an \( r \) interlacing random walk on \( A^N_n \) rather than \( A_n \). Similarly \( \psi_{u,v} \) is a bijection between \( \tau^N_{u,v} \) and \( \tau'^N_{u,v} \), and so Proposition \( \text{[3.5]} \) holds also in the discrete case.
We can extend the above constructions, by Kolmogorov consistency, to define a Markov chain \( \{ X^{(m)}(k), m \in \mathbb{Z} \} \) on the state space

\[
\{ x \in A_n^Z : x^{(m)} \preceq_r x^{(m+1)}, m \in \mathbb{Z} \}
\]

with the following properties:

1. For each \( m \in \mathbb{Z} \), \( (X^{(m)}(k), k \in \mathbb{Z}) \) is an \( s \)-interlacing random walk
2. For each \( k \in \mathbb{Z} \), \( (X^{(m)}(k), m \in \mathbb{Z}) \) is an \( r \)-interlacing random walk
3. For each \( m, k \in \mathbb{Z} \),

\[
X^{(m)}(k + 1) = \phi X^{(m+1)}(k), X^{(m)}(k+1) \}
\]

In the above, we can replace \( \phi \) by \( \psi \), and also consider the discrete versions. It is interesting to remark on the non-colliding random walk case, that is, the discrete model on \( A_n^N \) with \( r = s = 1/2\pi N \). In this case, the evolution of the above Markov chain is completely deterministic. Indeed, by ergodicity of the non-colliding random walk, for each \( k \in \mathbb{Z} \), with probability one, there exist infinitely many \( m \) such that \( X^{(m)}(k) \in \{(1 + c)/2\pi N, \ldots, (n + c)/2\pi N\}, c \in \mathbb{Z} \) and for these \( m \) there is only one allowable transition, namely to \( X^{(m)}(k + 1) = X^{(m)}(k) \) for \( j \leq n - 1 \) and \( X^{(m)}(k + 1) = X^{(m)}(k) + 1/2\pi N \). By property (3) above, this determines \( X^{(m)}(k + 1) \) for all \( m \in \mathbb{Z} \).

4. Couplings of interlaced Brownian motions

In this section we describe a Brownian motion construction that can considered as a scaling limit of the coupled random walks of the previous section.

The Laplacian on \( A_n \) with Dirichlet boundary conditions admits a unique non-negative eigenfunction, the function \( h \) defined previously at (12) which corresponds to the greatest eigenvalue \( \lambda_0 = -n(n-1)(n+1)/12 \). A \( h \)-Brownian motion on \( A_n \), is a diffusion with transition densities \( Q_t \) given by

\[
Q_t(x, x') = e^{-\lambda_0 \frac{h(x')}{h(x)}} Q^0_t(x, x'),
\]

where \( Q^0_t \) are the transition densities of standard \( n \)-dimensional Brownian motion killed on exiting \( A_n \), given explicitly by a continuous analogue of the Gessel-Zeilberger formula. [17, 21]. If \( (X(k); k \geq 0) \) is an \( h \)-Brownian motion in \( A_n \), then \( (e^{iX_1(t)}, e^{iX_2(t)}, \ldots, e^{iX_n(t)}) \) gives a process of \( n \) non-colliding Brownian motions on
the circle, the same as arises as the eigenvalue process of Brownian motion in the unitary group, see [12].

We wish to construct a bivariate process \((X, Y)\) with each of \(X\) and \(Y\) distributed as 
\(h\)-Brownian motions in \(A_n\), and with \(Y(t) \preceq_s X(t)\), for some fixed \(s \in (0, 2\pi)\). The dynamics we have in mind is based on the map \(\phi\) of the previous section. \(Y\) should be deterministically constructed from its starting point and the trajectory of \(X\), with \(Y_i(t)\) tracking \(X_i(t)\) except when this would cause the interlacing constraint to be broken. In the following we set \(\theta_0(t) = \theta_n(t)\).

**Proposition 4.1.** Given a continuous path \((x(t); t \geq 0)\) in \(A_n\), and a point \(y(0) \in A_n\) such that \(y(0) \preceq_s x(0)\), there exist unique continuous paths \((y(t); t \geq 0)\) in \(A_n\) and \((\theta(t); t \geq 0)\) in \(\mathbb{R}^n\) such that

(i) \(y_i(t) \preceq_s x_i(t)\) for all \(t \geq 0\);
(ii) \(y_i(t) - y_i(0) = x_i(t) - x_i(0) + \theta_{i-1}(t) - \theta_i(t)\);
(iii) for \(i = 1, 2, \ldots n\), the real-valued process \(\theta_i(t)\) starts from \(\theta_i(0) = 0\), is increasing, and the measure \(d\theta_i(t)\) is supported on the set \(\{t : y_{i+1}(t) = x_i(t)\}\).

Our principal tool in proving this proposition is another variant of the Skorohod reflection lemma. For \(l > 0\), let \(A_n(l) = \{x \in \mathbb{R}^n : x_1 \leq x_2 \leq \ldots \leq x_n \leq x_1 + l\}\). We adapt our usual convention, letting \(x_{n+1} = x_1 + l\) and \(x_0 = x_n - l\).

**Lemma 4.2.** Given a continuous path \((u(t); t \geq 0)\) in \(\mathbb{R}^n\) starting from \(u(0) = 0\), and a point \(v(0) \in A_n(l)\) then there exists a unique pair of \(\mathbb{R}^n\)-valued, continuous paths \((v, \theta)\) starting from \((v(0), 0)\) with

(i) \(v(t) \in A_n(l)\) for every \(t \geq 0\),
(ii) each component \(\theta_i(t)\) increasing with the measure \(d\theta_i(t)\) supported on the set \(\{t : v_{i+1}(t) = v_i(t)\}\),
(iii) \(v_i(t) = v_i(0) + u_i(t) - \theta_i(t)\).

**Proof.** Fix \(T > 0\). The path \(u\) is uniformly continuous on \([0, T]\), and so there exists \(\epsilon > 0\) such that \(\sum_i |u_i(t) - u_i(s)| < l/n^2\) for all \(|s - t| < \epsilon\). We prove existence and uniqueness over the time interval \([0, \epsilon]\), and then repeat the argument over consecutive time intervals \([\epsilon, 2\epsilon]\) etc.
Observe that there must exist $k \in \{1, 2, \ldots, n\}$ such that $v_{k+1}(0) - v_k(0) \geq t/n$. Set $v_k(t) = v_k(0) + u_k(t)$ and $\theta_k(t) = 0$ for $t \in [0, \epsilon]$. Then we follow the usual Skorohod lemma applied successively to give $v_{k-1}, v_{k-2}, \ldots, v_1, v_n, \ldots, v_{k+1}$, in particular we let

$$\theta_i(t) = \sup_{s \leq t} \left( v_{i+1}(t) - u_i(t) - v_i(0) \right)^-. $$

The choice of $\epsilon$ is such that the $v_{k-1}$ so constructed satisfies $v_k(t) < v_{k+1}(t)$ for all $t \in [0, \epsilon]$, and consequently all the desired properties of $v$ and $\theta$ hold. For uniqueness, we note that any $(v, \theta)$ for which the desired properties hold must be equal to the one just constructed, which follows from the uniqueness for the usual Skorohod construction.

**Proof of Proposition 4.1.** This is based on applying the Skorohod mapping to the path $u$ specified from $x$ by $u_i(t) = x_i(t) - x_i(0)$.

Choose $v(0) \in A_n(2\pi - s)$ so that $v_i(0) - v_{i-1}(0) = y_i(0) - x_{i-1}(0)$, and let $(v, \theta)$ be given by the Skorohod mapping for the domain $A_n(2\pi - s)$ with data $u$ and $v(0)$. Then for $t \geq 0$ define

$$y_i(t) = y_i(0) + u_i(t) + \theta_{i-1}(t) - \theta_i(t).$$

We claim that

(i) $y(t) \in A_n$ with $y(t) \preceq_x x(t)$ for every $t \geq 0$,

(ii) the measure $d\theta_i(t)$ is supported on the set \{t : $y_{i+1}(t) = x_i(t)$\}.

Calculate as follows.

$$y_i(t) - x_{i-1}(t) = (y_i(0) + u_i(t) + \theta_{i-1}(t) - \theta_i(t)) - (x_{i-1}(0) + u_{i-1}(t)) = (v_i(0) + u_i(t) - \theta_i(t)) - (v_{i-1}(0) + u_{i-1}(t) - \theta_{i-1}(t)) = v_i(t) - v_{i-1}(t).$$

This is valid even if $i = 0$ when we adhere to our conventions regarding $v_0$ etc. Assertion (ii) above follows since we see that \{t : $y_i(t) = x_{i-1}(t)$\} = \{t : $v_i(t) = v_{i-1}(t)$\}, and this latter set carries $d\theta_{i-1}$. Also since $v_i(t) \geq v_{i-1}(t)$ we obtain one part of the interlacing condition, namely, $y_i(t) \geq x_{i-1}(t)$. For the other part, consider the equality

$$x_i(t) - y_i(t) = x_i(0) - y_i(0) + \theta_i(t) - \theta_{i-1}(t).$$
Since this quantity is initially \( x_i(0) - y_i(0) > 0 \) and decreases only when \( \theta_i \) increases it follows that if there exists an instant \( t_1 \) for which \( x_i(t_1) - y_i(t_1) < 0 \), then there exists another instant, \( t_0 \), for which \( x_i(t_0) - y_i(t_0) < 0 \) and which belongs to the support of \( d\theta_i \). The latter implies that \( y_i(t_0) = x_{i-1}(t_0) \), and thus \( x_i(t_0) < y_i(t_0) = x_{i-1}(t_0) \) which contradicts \( x(t_0) \in A_n \). This proves existence. Uniqueness follows from the uniqueness statement in Proposition 4.2.

By virtue of this result we may make the following definition. For \( y(0) \in A_n \), let \( \Gamma_{y(0)} \) be the application which applied to an \( A_n \)-valued path \( (x(t); t \geq 0) \) returns the path \( (y(t); t \geq 0) \) specified by Proposition 4.1. The main result of this section is the following.

**Proposition 4.3.** Let \( (X(t); t \geq 0) \) be an \( h \)-Brownian motion in \( A_n \), starting from a point \( X(0) \) having the distribution \( q_{n}(y, dx) \) for some given \( y = y(0) \in A_n \). Then the process \( Y = \Gamma_{y(0)}(X) \) is distributed as a \( h \)-Brownian motion in \( A_n \).

The domain \( A_n \) is unbounded, which is a nuisance when we come to the probability, where life is easier if we have a finite invariant measure on the state space. Everything we do is invariant to shifts along the diagonal of \( \mathbb{R}^n \), and it is useful to project onto the hyperplane \( H_0 = \{ x \in \mathbb{R}^n : \sum x_i = 0 \} \). For \( x \in \mathbb{R}^n \), the orthogonal projection of \( x \) onto \( H_0 \) is given by \( x \mapsto x - \bar{x}1 \) where \( 1 \in \mathbb{R}^n \) is the vector with every component equal to 1, and \( \bar{x} = n^{-1} \sum x_i \). If \( (X(t); t \geq 0) \) is an \( h \)-Brownian motion in \( A_n \), then its projection onto \( H_0 \) is itself a diffusion process and indeed can be described as the \( h \)-transform of an \( (n - 1) \)-dimensional Brownian motion in \( H_0 \), killed on exiting \( A_n \cap H_0 \). Introducing more generally \( H_s = \{ x \in \mathbb{R}^n : \sum x_i = s \} \), and the notion of an \( h \)-Brownian motion on \( A_n \cap H_s \), we have the following variant of Proposition 4.3.

**Proposition 4.4.** Let \( (X(t); t \geq 0) \) be an \( h \)-Brownian motion in \( A_n \cap H_{s/2} \), starting from a point \( X(0) \) having the distribution \( q_{n}(y, dx) \) for some given \( y = y(0) \in A_n \cap H_{-s/2} \). Then the process \( Y = \Gamma_{y(0)}(X) \) is distributed as a \( h \)-Brownian motion in \( A_n \cap H_{-s/2} \).

Proposition 4.4 is easily deduced from this variant using the fact that if \( (X(t); t \geq 0) \) is an \( h \)-Brownian motion in \( A_n \), then the projection of \( X \) onto \( H_0 \) is independent of the process \( n^{-1/2} \sum X_i(t) \), which is a one-dimensional Brownian motion.
The results in the previous section were proved using the measure-preserving
properties associated with the dynamics for a single update. Since here we are work-
ing with continuous time processes such one time-step methods are not applicable.
The use of the Skorohod lemma in the construction of $\Gamma$ suggests that there may be
a role to be played by a certain reflected Brownian motion. We describe next how,
adapting the idea of Proposition 4.1 slightly, we can construct interlaced processes
$X$ and $Y$ from a reflected Brownian motion $R$ in the domain $H_0 \cap A_n(2\pi - s)$. Then
it will turn out that time reversal properties of $R$ can be used to prove Proposition
4.4.

Let $E(s) = \{(x, y) \in (A_n \cap H_{s/2}) \times (A_n \cap H_{-s/2}) : y \preceq_s x\}$. We can introduce
new co-ordinates on $E(s)$ as follows. For $(x, y) \in E(s)$ we let $f(x, y)$ be the unique
$(r, l) \in H_0 \times H_0$ such that

\begin{align}
  r_{i+1} - r_i &= y_{i+1} - x_i, \\
  l_{i+1} - l_i &= x_{i+1} - y_{i+1},
\end{align}

where $r_{n+1} = r_1 + (2\pi - s)$ and $l_{n+1} = l_1 + s$. It is easily seen that $f : E(s) \rightarrow
(H_0 \cap A_n(2\pi - s)) \times (H_0 \cap A_n(s))$ is bijective.

Now we construct a process in $E(s)$ via these alternative co-ordinates. Begin by
letting $(U(t); t \geq 0)$ be a standard Brownian motion in $\mathbb{R}^n$ starting from zero, and
let $V(0)$ be an independent random variable, uniformly distributed in $H_0 \cap A_n(2\pi - s)$. Let $(V, \Theta)$ be determined from $U$ and $V(0)$ by applying the Skorohod mapping
for $A_n(2\pi - s)$ as given in Proposition 4.2. Finally let $R(t)$ be the projection of $V(t)$
onto $H_0$, so $R(t) = V(t) - \bar{V}(t)1$. The process $(R(t); t \geq 0)$ is, by construction, a
semimartingale reflecting Brownian motion in the polyhedron $A_n(2\pi - s) \cap H_0$. See
41 for the general theory of such processes. Next we introduce a process $L$ also
taking value in $H_0$, which is constructed out a random initial value $L(0)$ together
with the increasing processes $\Theta_i$ for $i = 1, 2, \ldots, n$. Choose $L(0)$ independent of $R$
and uniformly distributed on $A_n(s) \cap H_0$, and let $L(t)$ be given by

$$
L(t) = L_i(0) - \pi \Theta(t),
$$

where $\pi : \mathbb{R}^n \rightarrow H_0$ denotes the projection onto $H_0$ defined by $\pi x = x - \bar{x}1$. Define
the stopping time $\tau = \inf\{t \geq 0 : L(t) \not\in A_n(s)\}$. Then for $t \leq \tau$, we define $X(t)$
and $Y(t)$ by $(R(t), L(t)) = f(X(t), Y(t))$. The joint law of $(X(t \wedge \tau), Y(t \wedge \tau); t \geq 0)$ may be described as follows.

1. $(X(0), Y(0))$ is uniformly distributed on $E(s)$.
2. $(X(t \wedge \tau); t \geq 0)$ is distributed as a Brownian motion in $H_0$ stopped at the instance it first leaves $A_n$, and conditionally independent of $Y(0)$ given $X(0)$.
3. $Y = \Gamma Y(0) X$.

We now turn to the time reversibility of $R$. For any vector $x \in \mathbb{R}^n$ we will denote by $x^\dagger$ the vector given by $x^\dagger_i = -x_{n-i+1}$. Note that if $x \in H_0 \cap A_n(2\pi - s)$ then $x^\dagger \in H_0 \cap A_n(2\pi - s)$ also. We also define, for any $x \in \mathbb{R}^n$, the vector $x^\ddagger$ via $x^\ddagger_i = -x_{n-i} - s/n$ for $1 = 1, 2, \ldots (n-1)$ and $x^\ddagger_n = -x_n + s(n-1)/n$. Note that if $x \in H_0 \cap A_n(s)$ then $x^\ddagger \in H_0 \cap A_n(s)$ also.

**Proposition 4.5.** Fix some constant $T > 0$. Let the processes $R$ and $L$ be as above and let $\Lambda$ be the event \{L(t) $\in A_n(s)$ for all $t \in [0, T]$\}. Then conditionally on $\Lambda$,

$$(R(t), L(t); t \in [0, T]) \overset{law}{=} (R^\dagger(T - t), L^\ddagger(T - t); t \in [0, T]).$$

**Proof.** As remarked above $R$ is a semimartingale reflected Brownian motion in the polyhedral domain $A_n(2\pi - s) \cap H_0$. Indeed it satisfies

$$R(t) = R(0) + B(t) + \frac{1}{\sqrt{2}} \sum \Theta_i(t)v^i,$$

where $B$ is a standard Brownian motion in $H_0$ and the vector $v^i$ describes the direction of reflection associated with the face $F_i = \{x \in A_n(2\pi - s) \cap H_0 : x_i = x_{i+1}\}$. Let $n^i$ be the inward facing unit normal to this face. Then an easy calculation shows that $v^i$, which is normalized so that the inner product $n^i \cdot v^i = 1$, is given by $v^i = n^i + q^i$ where the $j$th component of the vector $q^i$ is given by

$$q^i_j = \begin{cases} \sqrt{2}/n - 1/\sqrt{2} & \text{if } j = i, i + 1, \\ \sqrt{2}/n & \text{otherwise.} \end{cases}$$

We observe that the skew-symmetry condition,

$$n^i \cdot q^j + q^i \cdot n^j = 0,$$
for all \(i \neq j\), is met. Consequently by Theorem 1.2 of [41], the reflected Brownian motion \(R\) is in duality relative to Lebesgue measure to another reflected Brownian motion on \(H_0 \cap A_n(2\pi - s)\) with direction of reflection from the face \(F_i\) being \(n^i - q^i\). It is not difficult to check that \(R^i\) is such a reflected Brownian motion. Thus

\[
(R(t); t \in [0, T]) \overset{law}{=} (R^i(t - T); t \in [0, T]).
\]

The process \(2\Theta_i(t)\) is the local time of \(R_{i+1}(t) - R_i(t)\) at zero, and can be represented as

\[
\lim_{\epsilon \downarrow 0} \frac{1}{\epsilon} \int_0^t 1(R_{i+1}(s) - R_i(s) \leq \epsilon)ds.
\]

Now note that,

\[
\int_0^t 1(R_{i+1}(s) - R_i(s) \leq \epsilon)ds \overset{law}{=} \int_0^t 1(R^i_{i+1}(T - s) - R^i_i(T - s) \leq \epsilon)ds = \int_{T-t}^T 1(R_{n-1+1}(s) - R_{n-1}(s) \leq \epsilon)ds.
\]

From this we deduce that the time reversal property extends to

\[
(R(t), \Theta(t); t \in [0, T]) \overset{law}{=} (R^i(t - T), \Theta^i(T) - \Theta^i(T - t); t \in [0, T]),
\]

where \(\Theta^i_i(t) = \Theta_{n-i}\) for \(i = 1, 2, \ldots, (n - 1)\) and \(\Theta^i_n = \Theta_n\).

Let \(F\) be a bounded path functional. Let \(v_n(s)\) denote the Lebesgue measure of \(H_0 \cap A_n(s)\). Then using the time reversal property,

\[
\mathbb{E}[F(R(t), L(t); t \in [0, T])1_A] = \\
\frac{1}{v_n(s)} \mathbb{E} \left[ \int_{H_0} d\alpha F(R(t), \alpha - \pi \Theta(t); t \in [0, T])1\{\alpha - \pi \Theta(t) \in A_n(s); t \in [0, T]\} \right] = \\
\frac{1}{v_n(s)} \mathbb{E} \left[ \int_{H_0} d\alpha F(R^i(t - T), \alpha - \pi \Theta^i(T) + \pi \Theta^i(T - t); t \in [0, T]) \right. \\
\times 1\{\alpha - \pi \Theta^i(T) + \pi \Theta^i(T - t) \in A_n(s); t \in [0, T]\}. 
\]
Now we make the substitution $\hat{\alpha} = \alpha - \pi \Theta^\dagger(T)$ to obtain

$$\frac{1}{v_n(s)} \mathbb{E} \left[ \int_{H_0} \! d\hat{\alpha} F(R^\dagger(T - t), \hat{\alpha} + \pi \Theta^\dagger(T - t); t \in [0, T]) \times 1\{\hat{\alpha} + \pi \Theta^\dagger(T - t) \in A_n(s), t \in [0, T]\} \right] = \mathbb{E} \left[ F(R^\dagger(T - t), L^\dagger(T - t); t \in [0, T]) 1_A \right],$$

where we have used $L^\dagger(t) = L^\dagger(0) + \pi \Theta^\dagger(t)$. □

**Proof of Proposition 4.4.** Let $R$, $\Theta$ and $L$ be as above, and once again let $A$ be the event that $\{L(t) \in A_n(s) \text{ for all } t \in [0, T]\}$. Recall the mapping $f$ such that $(R(t), L(t)) = f(X(t), Y(t))$. It is easily verified that $(Y^\dagger(t), X^\dagger(t)) \in E(s)$ and that

$$(R^\dagger(t), L^\dagger(t)) = f(Y^\dagger(t), X^\dagger(t)).$$

Thus the preceding time reversal result implies that, conditionally on $A$,

$$(X(t), Y(t); t \in [0, T]) \overset{\text{law}}{=} (Y^\dagger(T - t), X^\dagger(T - t); t \in [0, T]).$$

For the final step of the argument we consider $X$ and $Y$ be as above and denote the governing measure by $P$. Then we let

$$\tilde{P} = e^{-\lambda_0 T} \gamma T \mathbb{1}_A h(Y(0))h(X(T)) \cdot P.$$

Under $\tilde{P}$, the equality in law,

$$(X(t), Y(t); t \in [0, T]) \overset{\text{law}}{=} (Y^\dagger(T - t), X^\dagger(T - t); t \in [0, T]).$$

holds unconditionally. Finally we note that under $\tilde{P}$, the distribution of $(X(t), t \in [0, T])$ and hence of $(Y^\dagger(T - t), t \in [0, T])$ is that of a stationary $h$-Brownian motion on $A_n \cap H_{s/2}$. But this latter law is invariant under time reversal, and its image under the conjugation $x \mapsto x^\dagger$ is the law of a stationary $h$-Brownian motion on $A_n \cap H_{-s/2}$. This is therefore the law of $(Y(t), t \in [0, T])$. Conditioning on $Y(0)$ gives the statement of the proposition. □

In the case $n = 2$, the results of this section can be expressed in terms of Brownian motion in a compact interval. Let $X = (X_t, t \geq 0)$ be a Brownian motion conditioned, in the sense of Doob, never to exit the interval $[-p, p]$, where $p > 0$. Let $y \in [-p, p]$, $a \in [0, 2p]$ and suppose that the initial law of $X$ is supported
in the interval \(|y + a - p| - p, p - |y + p - a|\) with density proportional to \(\cos(\pi x/2p)\). Let \(Z\) be the image of the path \(X + (y - X_0 + a)/2\) under the Skorohod reflection map for the interval \([0, a]\). In other words,

\[
Z_t = X_t + (y - X_0 + a)/2 + L_t - U_t,
\]

where \(L\) and \(U\) are the unique continuous, non-decreasing paths such that the points of increase of \(L\) occur only at times when \(Z_t = 0\), the points of increase of \(U\) occur only at times when \(Z_t = a\), and \(Z_t \in [0, a]\) for all \(t \geq 0\). Then the process

\[
Y_t = y - X_0 + X_t + 2(L_t - U_t) \quad t \geq 0
\]

is a Brownian motion conditioned, in the sense of Doob, never to exit the interval \([-p, p]\). This is a special case of Proposition 4.4. Actually, in the statement of that proposition we have \(p = \pi\), but this can be easily modified for general \(p\). It is interesting to consider this statement when \(y = 0\) and \(p \to \infty\). Then \(X\) is a standard Brownian motion, initially uniformly distributed on the interval \([-a, a]\). The process \(Z\) is a reflected Brownian motion in \([0, a]\), initially uniformly distributed on \([0, a]\). The conclusion in this case is that \(Y\) is a standard Brownian motion started from zero. We remark that in this setting, if instead we take \(X_0 = -a\), then \(Y\) is a Brownian motion started from zero, conditioned (in an appropriate sense) to hit \(a\) before returning to zero. This is a straightforward consequence of the above result (for uniform initial law) and the fact (see [37]) that, if we set \(T = \inf\{t \geq 0 : Y_t = a\}\), then the law of \(X_T\) is uniform on \([-a, a]\). Note that if we let \(a \to \infty\) in this case we recover Pitman’s representation for the three-dimensional Bessel process. There are explicit formulae for the Skorohod reflection map for the compact interval \([0, a]\) and hence for the process \(Y\) in the above discussion. Let \(f(t) = X_t + (y - X_0 + a)/2\) and write \(f(s, t) = f(t) - f(s)\). A discrete version of the Skorohod problem was considered in [39], from which we deduce the expressions

\[
Z_t = \max \left\{ \sup_{0 \leq r \leq t} \min\{f(r, t), a + \inf_{r \leq s \leq t} f(s, t)\}, \min\{f(t), a + \inf_{0 < s < t} f(s, t)\} \right\}
\]

\[
= \min \left\{ \inf_{0 \leq r \leq t} \max\{a + f(r, t), \sup_{r \leq s \leq t} f(s, t)\}, \max\{f(t), \sup_{0 < s < t} f(s, t)\} \right\}.
\]

An alternative formula was obtained in [26], which yields

\[
Z_t = \phi(t) - \sup_{0 \leq s \leq t} \left[ (\phi(s) - a)^+ \wedge \inf_{s \leq u \leq t} \phi(u) \right]
\]
where
\[ \phi(t) = f(t) + \sup_{0 \leq s \leq t} [-f(s)]^{+}. \]
It could be interesting to relate the corresponding expressions for \( Y_t \) to the Pitman transforms introduced in [4].

5. A bead model on the cylinder

In this section it will be convenient to work with a slightly weaker notion of interlacing, defined as follows. For \( a, b \in D_n \), write \( a \prec b \) if
\[ a_1 \leq b_1 < a_2 \leq \cdots < a_n \leq b_n, \]
and \( a \succ b \) if
\[ b_1 < a_1 \leq b_2 < \cdots \leq b_n < a_n. \]
For \( y = \{e^{ia_1}, \ldots, e^{ia_n}\} \) and \( z = \{e^{ib_1}, \ldots, e^{ib_n}\} \), where \( a, b \in D_n \), write \( y \prec z \) if either \( a \prec b \) or \( a \succ b \), and define
\[
l(y, z) = \begin{cases} \sum_j (b_j - a_j) & \text{if } \sum_j (b_j - a_j) \geq 0 \\ \sum_j (b_j - a_j) + 2\pi & \text{otherwise}. \end{cases}
\]
Consider the Markov kernels defined, for \( q > 0 \), by
\[
m_q(y, dz) = c_q^{-1} \int_0^{2\pi} |1 - e^{i r}|^{n-1} q^r p_r(y, dz)dr,
\]
where \( p_r \) is defined by (9) and
\[
c_q = \int_0^{2\pi} |1 - e^{i r}|^{n-1} q^r dr.
\]
By Proposition 2.1 if we define
\[
I_q(y, z) = \begin{cases} q^{l(y, z)} & \text{if } y \prec z, \\ 0 & \text{otherwise}, \end{cases}
\]
then
\[
m_q(y, dz) = \tilde{c}_q^{-1} \frac{\Delta(z)}{\Delta(y)} I_q(y, z)dz,
\]
where \( \tilde{c}_q = c_q/(n-1)! \). Recall that \( \mu(dx) = (2\pi)^{-n} \Delta(x)^2 dx \) is the probability measure on \( C_n \) induced from Haar measure on \( U(n) \). The Markov chain with
transition density $m_q$ has $\mu$ as an invariant measure and, with respect to $\mu$, has time-reversed transition probabilities

$$ m_q(z, dy) = \tilde{c}^{-1} \frac{\Delta(y)}{\Delta(z)} \mathcal{I}_q(y, z) dy. $$

We can thus construct a two-sided stationary version of this Markov chain to obtain a probability measure $\alpha$ on $C_n^Z$, supported on configurations $\cdots \prec x^1 \prec x^0 \prec x^1 \prec x^2 \prec \cdots$. We will show that $\alpha$ defines a determinantal point process on $[0, 2\pi]^Z$. By stationarity it suffices to consider the restrictions $\alpha_m$ to the cylinder sets $C_{n,m} := C_{n}^{\{1,2,\ldots,m\}}$. Writing $\bar{x} = (x^1, \ldots, x^m)$ and $d\bar{x} = dx^1 \cdots dx^m$,

$$ \alpha_m(d\bar{x}) = \mu(dx^1)m_q(x^1, dx^2) \cdots m_q(x^{m-1}, dx^m). \tag{25} $$

Assume for the moment that $q \neq 1$. Define a function $f : \mathbb{R} \to \mathbb{C}$ by

$$ f(u) = \frac{(qe^{i(n-1)/2})^u \mod 2\pi}{1 - (-1)^{n-1}q^{2\pi}}. $$

**Lemma 5.1.** For $y = \{e^{ia}, \ldots, e^{ia_n}\}$ and $z = \{e^{ib}, \ldots, e^{ib_n}\}$, where $a, b \in D_n$,

$$ \mathcal{I}_q(y, z) = (1 - (-1)^{n-1}q^{2\pi})e^{i\sum_{1\leq j,k \leq n}(a_j - b_j)} \det (f(b_k - a_j))_{1\leq j,k \leq n}. $$

**Proof.** Let $c \neq 1$ and consider the $n \times n$ matrix $W = (w_{jk})$ defined by

$$ w_{jk} = \begin{cases} 
1 & a_j \leq b_k \\
\left(\begin{array}{c}
\alpha_j \\
\eta
\end{array}\right) & a_j > b_k.
\end{cases} $$

If $a \prec b$, $W$ consists of 1’s on and above the diagonal and $c$’s below, so that $\det W = (1 - c)^{n-1}$. If $a \succ b$, $W$ consists of 1’s above the diagonal and $c$’s on and below the diagonal, so that $\det W = c(c - 1)^{n-1}$. If neither $a \prec b$ or $a \succ b$, then there must exist an index $j$ such that, either $a_j = a_{j+1}$, or $a_j < a_{j+1} \leq b_k$ for all $k$, or $b_k < a_j < a_{j+1} \leq b_{k+1}$ for some $k$. In each of these cases, rows $j$ and $j + 1$ of $W$ are identical and hence $\det W = 0$. Thus,

$$ \det W = \begin{cases} 
(1 - c)^{n-1} & a \prec b \\
c(c - 1)^{n-1} & a \succ b \\
0 & \text{otherwise.}
\end{cases} \tag{26} $$
Taking \( c = (q e^{i(n-1)/2} \pi) = (-1)^{n-1} q^{2 \pi} \), we can write

\[
(1 - c) e^{\frac{1}{2} \sum_{j} (a_j - b_j)} \det(f(b_j - a_k))_{1 \leq j, k \leq n} = q^{\sum_{j} (b_j - a_j)} (1 - c)^{-(n-1)} \det W
\]

\[
= \begin{cases} 
q^{\sum_{j} (b_j - a_j)} & a \prec b \\
q^{\sum_{j} (b_j - a_j) + 2\pi} & a \succ b \\
0 & \text{otherwise}
\end{cases}
\]

\[
= \mathcal{I}_q(y,z),
\]

as required. \( \square \)

For \( r = 1, \ldots, m - 1 \), define \( \phi_{r,r+1} : [0, 2\pi)^2 \rightarrow \mathbb{C} \) by \( \phi_{r,r+1}(a,b) = f(b-a) \).

Define \( \phi_{0,1} : \mathbb{R} \times [0, 2\pi) \rightarrow \mathbb{C} \) and \( \phi_{m,m+1} : [0, 2\pi) \times \mathbb{R} \rightarrow \mathbb{C} \) by \( \phi_{0,1}(a,b) = e^{ia} \)

and \( \phi_{m,m+1}(a,b) = e^{-ia} \). For \( r = 1, \ldots, m \), write \( x^r = \{ e^{ia_1}, \ldots, e^{ia_n} \} \), where \( a^r \in D_n \), and set \( a_j^0 = a_j^{m+1} = j - 1 \), for \( j = 1, \ldots, n \).

**Theorem 5.2.** For \( q \neq 1 \),

\[
\alpha_m(d\vec{x}) = Z_m^{-1} \prod_{r=0}^{m} \det \left( \phi_{r,r+1}(a_j^r, a_k^{r+1}) \right)_{1 \leq j, k \leq n} d\vec{x},
\]

where \( Z_m = \tilde{c}_q^{m-1} (1 - (-1)^{n-1} q^{2\pi} )^{-(m-1)} (2\pi)^n \).

**Proof.** By (24) we can write

\[
\alpha_m(d\vec{x}) = \tilde{c}_q^{-(m-1)} (2\pi)^{-n} \Delta(x^1) \Delta(x^m) \mathcal{I}_q(x^1, x^2) \cdots \mathcal{I}_q(x^{n-1}, x^m) d\vec{x}.
\]

Using the formula

\[
\Delta(x^1) \Delta(x^m) = \det \left( e^{i(j-(n+1)/2)a^1_k} \right)_{1 \leq j, k \leq n} \det \left( e^{-i(j-(n+1)/2)a^n_j} \right)_{1 \leq j, k \leq n},
\]
and Lemma 5.1 we obtain
\[\alpha_m(d\bar{x}) = Z_m^{-1} e^{\frac{2\pi i}{m+1} \sum_j (a_j^m - a_j^m)} \det \left( e^{i(j-(n+1)/2)a_j^m} \right)_{1 \leq j, k \leq n} \times \det \left( e^{-i(j-(n+1)/2)a_k^{m+1}} \right)_{1 \leq j, k \leq n} \prod_{r=1}^{m-1} \det \left( \phi_{r,r+1}(a_j^r, a_k^{r+1}) \right)_{1 \leq j, k \leq n} d\bar{x} = Z_m^{-1} \det \left( e^{i(j-1)a_k^1} \right)_{1 \leq j, k \leq n} \det \left( e^{-i(j-1)a_k^{n+1}} \right)_{1 \leq j, k \leq n} \times \prod_{r=1}^{m-1} \det \left( \phi_{r,r+1}(a_j^r, a_k^{r+1}) \right)_{1 \leq j, k \leq n} d\bar{x} = Z_m^{-1} \prod_{r=0}^{m-1} \det \left( \phi_{r,r+1}(a_j^r, a_k^{r+1}) \right)_{1 \leq j, k \leq n} d\bar{x},\]
as required. \( \square \)

**Corollary 5.3.** For any \( q > 0 \), the measure \( \alpha \) defines a determinantal point process on \([0, 2\pi)^2\) with space-time correlation kernel given by
\[
K(r, a; s, b) = \begin{cases} 
\frac{1}{2\pi} \sum_{k=0}^{n-1} g_k^{r-s} e^{i(b-a)k} & r \geq s \\
-\frac{1}{2\pi} \sum_{k \in \mathbb{Z} \setminus \{0, \ldots, n-1\}} g_k^{r-s} e^{i(b-a)k} & r < s 
\end{cases}
\]
where
\[
g_k = \left( \int_0^{2\pi} f(u) e^{-iku} du \right)^{-1} = i \left( k - \frac{n-1}{2} \right) - \log q.
\]

**Proof.** For \( q \neq 1 \), this follows from Theorem 5.2 [22 Proposition 2.13] and a straightforward computation. The case \( q = 1 \) is obtained by continuity. \( \square \)

Lemma 5.1 can be used to give a direct proof of [24], and hence Proposition 2.1

**Proof of Proposition 2.1.** The characters \( \chi_\lambda \) are given, for \( y = \{e^{ia_1}, \ldots, e^{ia_n}\} \), by
\[
\chi_\lambda(y) = i^{-\binom{n}{2}} \Delta(y)^{-1} \det (e^{i\mu_j \Delta_k})_{1 \leq j, k \leq n},
\]
where \( \mu = \lambda + \rho \) and
\[
\rho = \left( \frac{n-1}{2}, \ldots, \frac{n-1}{2} - 1 \right) = \left( \frac{n-1}{2}, \ldots, \frac{n-1}{2} + 1 \right).
\]
Using Lemma 5.1 and the Cauchy-Binet formula, we obtain
\[
\int \frac{\Delta(z)}{\Delta(y)} I_\phi(y, z) \chi_\lambda(z) dz = (1 - (-1)^{n-1} q^{2\pi}) \prod_j (-i\mu_j - \log q)^{-1} \chi_\lambda(y).
\]
On the other hand, writing \( x_r = \{e^{ir}, 1, 1, \ldots, 1 \} \), an easy calculation shows that
\[
\int_0^{2\pi} |1 - e^{ir}y| \frac{\chi_{\lambda}(x_r)}{d_{\lambda}} dr = (n - 1)! (1 - (-1)^{n-1}q^{2\pi}) \prod_j (-i\mu_j - \log q)^{-1}
\]
and so, by (7),
\[
m_q\chi_{\lambda} = \hat{c}^{-1}_q (1 - (-1)^{n-1}q^{2\pi}) \prod_j (-i\mu_j - \log q)^{-1}\chi_{\lambda}.
\]
Since \( \{\chi_{\lambda}, \lambda \in \Omega_n\} \) is a basis for \( L_2(C_n, \mu) \), this implies (24).

Analogous results to those presented in this section can be obtained for the discrete version of this model, which is equivalent to considering a certain family of Gibbs measures on rhombic tilings of the cylinder. For more details, see [29]. The couplings defined in Section 3 are quite useful in this setting, where the group-theoretic considerations of Section 2 no longer apply. For example, they can be used to prove that the discrete analogues of the interlacing operators \( \{I_q, q > 0\} \) commute with each other. In this setting, the symmetric functions
\[
(q_1, \ldots, q_k) \mapsto I_{q_1} \cdots I_{q_k}(y, z)
\]
are essentially the cylindrical skew Schur functions discussed in the papers [27, 36].

Finally, we remark that, in the case \( q = 1 \), the probability measure \( \alpha_2 \) defined by (25) also arises naturally in random matrix theory. The probability measures on \( C_n \) given by \( \mu(dx) = (2\pi)^{-n} \Delta(x)^2 dx \) and \( A_n^{-1} \Delta(x) dx \), where \( A_n \) is a normalisation constant, are known, respectively, as the circular unitary ensemble and circular orthogonal ensemble. It is a classical result, which was conjectured by Dyson [11] and subsequently proved by Gunson [20], that the set of alternate eigenvalues from a superposition of two independent draws from the circular orthogonal ensemble, are distributed according to the circular unitary ensemble. Moreover, the joint law of the ‘even’ and ‘odd’ eigenvalues has probability density on \( C_n \times C_n \) proportional to \( \Delta(y)\Delta(z)I_1(y, z) \), which is the same as \( \alpha_2 \), the joint distribution at two consecutive times of the stationary Markov chain with transition kernel \( m_q \) (defined by (23)) in the case \( q=1 \).

Acknowledgements. Research of the first two authors supported in part by Science Foundation Ireland Grant No. SFI04/RP1/I512. Thanks are due to Dominique Bakry, Manon Defosseux and Persi Diaconis for helpful discussions, and also to
the anonymous referee and Associate Editor for helpful comments and suggestions which have led to an improved version of the paper.

References

[1] D. Bakry and N. Huet. The hypergroup property and representation of Markov kernels. In: Séminaire de Probabilités XLI, Lecture Notes in Mathematics, Vol. 1934, Springer 2008.
[2] Y. Baryshnikov. GUEs and queues. Probab. Theory Related Fields, 119 (2001) 256-274.
[3] P.H. Berard. Spectres et groupes cristallographiques. I. Domaines euclidiens. Invent. Math. 58 (1980) 179.
[4] Ph. Biane, Ph. Bougerol, N. O’Connell. Littlewood paths and Brownian paths. Duke Math. J. 130 (2005) 127–167.
[5] C. Boutillier. The bead model and limit behaviors of dimer models. Ann. Probab., to appear.
[6] C. Boutillier and B. de Tilière. Loops statistics in the toroidal honeycomb dimer model. arXiv:math/0608600.
[7] Manon Defosseux. Orbit measures and interlaced determinantal point processes. arXiv:0802.4183.
[8] P. Diaconis. Patterns in eigenvalues: the 70th Josiah Williard Gibbs lecture Bull. Amer. Math. Soc. 40 (2003), no. 2, 155–178.
[9] P. Diaconis and M. Shahshahani. Products of random matrices as they arise in the study of random walks on groups. Contemp. Math. 50 (1986) 183–195.
[10] Persi Diaconis and James A. Fill. Strong stationary times via a new form of duality. Ann. Probab., 18:1483–1522, 1990.
[11] F. Dyson. Statistical theory of the energy levels of complex systems, I-III. J. Math. Phys. 3 (1962).
[12] F. Dyson. A Brownian-motion model for the eigenvalues of a random matrix. J. Math. Phys. 3 (1962).
[13] Jacques Faraut. Analysis on Lie Groups: An Introduction. Cambridge University Press, 2008.
[14] P.J. Forrester and T. Nagao. Determinantal correlations for classical projection processes. arXiv:0801.0100
[15] P. Forrester, E. Rains. Jacobians and rank 1 perturbations relating to unitary Hessenberg matrices. Int. Math. Res. Not., 2006.
[16] W. Fulton. Young Tableaux: With Applications to Representation Theory and Geometry. Cambridge University Press, 1997.
[17] I. Gessel, R. Zeilberger. Random walk in a Weyl chamber. Proc. Amer. Math. Soc. 115. (1992), 27-31.
[18] David J. Grabiner. Random Walk in an Alcove of an Affine Weyl Group, and Non-Colliding Random Walks on an Interval. J. Combin. Theor., Series A, (2002) 285–306.
[19] J. Gravner, C. Tracy, H. Widom. Limit theorems for height fluctuations in a class of discrete space and time growth models. *J. Stat. Phys.* 102 (2001), 1085-1132.

[20] J. Gunson. Proof of a Conjecture of Dyson in the Statistical Theory of Energy Levels. *J. Math. Phys.* 3 (1962) 752-753.

[21] D. Hobson and W. Werner. Non-Colliding Brownian Motions on the Circle *Bull. London Math. Soc.* 28 (1996) 643–650.

[22] Kurt Johansson. Random matrices and determinantal processes Lectures given at the summer school on Mathematical statistical mechanics in July 05 at Ecole de Physique, Les Houches. arXiv:math-ph/0510038.

[23] K. Johansson and E. Noordenstam. Eigenvalues of GUE minors. *Elect. J. Probab.* 11 (2006) 1342–1371.

[24] R. Kenyon, A. Okounkov, S. Sheffield. Dimers and amoebae. *Ann. Math.* 163, no. 3 (2006) 1019–1056.

[25] Wolfgang König, Neil O’Connell and Sebastian Roch. Non-colliding random walks, tandem queues, and discrete orthogonal polynomial ensembles. *Electron. J. Probab.* 7 (2002).

[26] Lukasz Kruk, John Lehoczky, Kavita Ramanan and Steven Shreve. An explicit formula for the Skorokhod map on [0,a]. *Ann. Probab.* 35 (2007), 1740-1768.

[27] P. McNamara Cylindric skew Schur functions. *Adv. Math.* 205 (2006), no. 1, 275–312.

[28] M.L. Mehta. *Random Matrices.* Third edition. Elsevier, 2004.

[29] Anthony P. Metcalfe. Phd Thesis, in preparation.

[30] N. O’Connell. A path-transformation for random walks and the Robinson-Schensted correspondence. *Trans. Amer. Math. Soc.* 355 (2003) 3669–3697.

[31] N. O’Connell. Conditioned random walks and the RSK correspondence. *J. Phys. A*, 36:3049–3066, 2003.

[32] N. O’Connell, M. Yor. A representation for non-colliding random walks. *Electron. Comm. Probab.* 7 (2002).

[33] A. Okounkov and N. Reshetikhin. The birth of random matrix. *Moscow Math. J.* 6 (2006) 553–566.

[34] U. Porod. The cut-off phenomenon for random reflections *Ann. Probab.* 24 (1996) 74–96.

[35] U. Porod. The out-of phenomenon for random reflections II: Complex and quaternionic cases. *Probab. Theory Related Fields* 104 (1996) 181–209.

[36] A. Postnikov Affine approach to quantum Schubert calculus. *Duke Math. J.,* 128 (2005) 473-509.

[37] L. C. G. Rogers and J. Pitman. Markov functions. *Ann. Probab.* 9:573–582, 1981.

[38] J.S. Rosenthal. Random rotations: characters and random walks on SO(N) . *Ann. Probab.* 22 (1994) 398–423.

[39] Fergal Toomey. Bursty traffic and finite capacity queues. *Ann. Oper. Res.* 79 (1998) 45-62.

[40] Jon Warren. Dyson’s Brownian motions, intertwining and interlacing. *Electron. J. Probab.* 12 (2007), no. 19, 573–590.
[41] R. J. Williams. Reflected Brownian motion with skew symmetric data in a polyhedral domain. 
Probab. Theory and Related Fields 75 (1987) 459–485.

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