Spectral domain inverse design for accelerating nanocomposite metamaterials discovery

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Abstract: Inverse design techniques in the context of nanophotonics have helped in discovery of compact and counter-intuitive structures/shapes. We introduce the concept of spectral domain inverse design to search through the optical trade-space (dispersive permittivity) of nanocomposite metamaterials. We develop a hybrid optimization technique that combines genetic algorithms and gradient descent methods. We utilize this technique to inverse design an ultra-thin thermophotovoltaic emitter coating material. Our work can lead to an efficient approach to search for new multi-functional optical/thermal metamaterials with desired complex permittivity.

1. Introduction

Topology optimization, machine learning, gradient descent, needle optimization and genetic algorithms have risen to the forefront of nanophotonics due to their ability to inverse design nanophotonic devices for specific applications [1–19]. Recent important examples include on-chip wavelength demultiplexer [3], multilayer thin films [4–9], chirped-mirrors [10], nanoparticle scattering [15], metasurfaces [16,18], core-shell nanoparticles [19] to name a few. The common theme in these works is the optimization of structural (topology/shape optimization) [3,14] or configurational (various possible combinations in multilayer thin films) [8,9] parameters for a given set of available materials to achieve predefined target goals. This motivates the search for analogous techniques that search through the optical dispersion (dispersive nature of material’s permittivity) trade-space to enable discovery of new nanocomposites and metamaterials.

Searching through optical material dispersion space is challenging. Our aim in this work is to put forth spectral domain inverse design as an approach to achieve desired optical responses through nanocomposite materials and metamaterials. The goal is to search for an optimum material platform with an engineered complex dielectric permittivity that depends on frequency. Thus the target functionality is achieved by changing the effective dielectric permittivity of nano-composites. Coupled with existing structural [3,14] (i.e. shape) or configurational optimization [8,9], our approach can open the route to utilize unique metal-dielectric-semiconductor nanocomposites as well as unconventional emerging 2D materials to realize desired optical responses.

We perform a proof of principle optimization to show that spectral nanophotonic design can lead to the discovery of new metamaterials in an efficient manner. Conventional plasmonic thin films of gold, silver, etc. form the starting point for nano-structuring into metamaterials. To discover alternative optimal nanocomposite thin films with desired broadband spectral properties, we utilize effective medium theory in the inverse design algorithm. In our technique, we make use of genetic algorithms and gradient descent techniques in an alternating optimization scheme-thus we call it “Hybrid Optimization”. We show that conventional multi-layer thin films can be replaced by an optimized metamaterial nanocomposite Fig 1. This hybrid optimization method...
is used to find an optimal single layer disordered metamaterial for thermal photovoltaic (TPV) emitter coating applications.

![Spectral features of a N layered multilayer thermal emitter (TPV emitter) stack are mimicked by a single disordered nanocomposite metamaterial. Using a hybrid optimization technique we find the effective permittivity of this nanocomposite material that as the same reflection and transmission spectral features as the one shown.](image)

Figure 1. **Basic Idea:** The spectral features of a N layered multilayer thermal emitter (TPV emitter) stack are mimicked by a single disordered nanocomposite metamaterial. Using a hybrid optimization technique we find the effective permittivity of this nanocomposite material that as the same reflection and transmission spectral features as the one shown.

## 2. Hybrid Optimization

To start with, here we outline how the hybrid optimization is implemented. First, we specify the desired reflection, transmission and absorption spectra (total of three spectra). In our approach, unlike other implementations [8, 9], our optimization is highly constrained as we simultaneously optimize on reflection $R^*(\lambda)$, and transmission $T^*(\lambda)$ target spectra. The spectral response of the nanocomposite metamaterial is computed via transfer matrix method. Effective Medium Theory (EMT) [20–23] is used to compute the effective permittivity of the nanocomposite metamaterial, as shown in Eq.1.

$$\epsilon_{eff} = \epsilon_{mat} - \frac{3\pi\epsilon_h}{2x^3} \left[ \frac{\Sigma_{l=1}^{N} i(2l+1)(a_l + b_l)}{1 + \frac{\rho\epsilon_h}{2x^3} \sum_{l=1}^{N} l(2l+1)(a_l + b_l)} \right]$$

(1)

where, $x = 2\pi r n_k/\lambda$ is the size parameter ($r$ is the radius, $\epsilon_h(n_k)$ is permittivity (refractive index) of host medium and $\lambda$ is the wavelength of light), $\epsilon_{mat}$ corresponds to the permittivity of the inclusion material (nanoparticle’s material), $l$ is the order of the poles, $a_l$ and $b_l$ are the Mie coefficients that are expressed in terms of Ricatti-Bessel functions. The size parameter sets the cut-off for the number of terms to be considered. It is defined as, $N = [x + 4.05x^\dagger + 2]$, where the square brackets mean rounding it to the nearest integer. This enables the equation Eq.1 above to handle effects arising due to finite size spherical particle inclusions. This effective medium theory provides the additional degree of freedom in material permittivity. Here, we use this idea to find an equivalent single layer nanocomposite metamaterial.

To find the optimum nanocomposite structure, we try to minimize the residuals between the specified target spectra as described by the merit function $MF(d, r, \rho, n^i, n^t)$. The merit function is defined as Eq.(2), where $d$ is the thickness of the single layer, $r$ is the radius of the nanoparticle inclusions, $\rho$ is the fill fraction, $n^i$ and $n^t$ are the material indices corresponding to a discrete material set. We have used a discrete material set to account for practical limitations. In this specific problem we have used a total of 22 materials.
Figure 2. **Hybrid Optimization:** The flowchart shows the algorithm used to obtain the optimal solution. Gradient descent technique was used for optimizing $d$, $\rho$, $r$, and a discrete search (Genetic algorithm) was used for searching over the material database.

\[
MF = \sum_{\lambda} \sqrt{\left[ (R(\lambda) - R^*(\lambda))^2 + (T(\lambda) - T^*(\lambda))^2 + (A(\lambda) - A^*(\lambda))^2 \right]},
\]

where $A^*(\lambda)$ is $1 - R^*(\lambda) - T^*(\lambda)$. As we are interested in optimizing over a broad wavelength range, we consider the material dispersion for accurate results. Given the desired target spectra, $R^*(\lambda)$ and $T^*(\lambda)$, the optimization problem is formulated as a mixed integer problem,

\[
[d^*, r^*, \rho^*, n_1^*, n_2^*] = \arg\min MF(d, r, \rho, n_1, n_2),
\]

where, $d, d^*, r, r^*, \rho, \rho^* \in \mathbb{R}^+$ and $n_1, n_1^*, n_2, n_2^* \in \mathbb{N}$.

Having defined the optimization problem in Eq.(2), we now describe the hybrid optimization approach that we have implemented. One can classify this optimization as a non-convex optimization problem. The merit function space is very complex and simultaneous optimization of all the parameters did not yield a satisfactory solution (the optimization often converged to a shallow local minimum). As part of the problem considered here has discrete variables, the problem falls in a class of combinatorial optimization. Many problems in this class are NP-hard and efficient optimal algorithms are infeasible unless $P=NP$ [24]. NP-hard are class of problems which are at least as hard as the hardest problems in NP where NP is defined as a class of computational problems where given an answer, one can verify the solution in polynomial time by a deterministic Turing machine.

Considering the complexity of the problem, we have utilized hybrid optimization approach wherein we use a coordinate descent method. A schematic describing the basic algorithmic idea is shown in Fig.2. In this method, we optimize one parameter keeping the rest of the parameters fixed. Once the favourable solution for the first parameter is optimized, the next (second) parameter is now optimized using the value of the first optimal parameter as the new initial guess value for the second parameter’s optimization run. This process is iteratively repeated across all the parameters until the desired convergence criteria is satisfied. Since each sub-problem of optimizing one parameter at a time converges (decreasing) and the overall problem is bounded from below, the hybrid algorithm converges to a stationary solution [25, 26].

The hybrid approach can primarily be broken down into the following steps: (A) Initialization, (B) Optimal gradient descent, and (C) A discrete optimization. Considering the non-convex nature of the merit function space, it is very important to have an intelligent guess of the initialization parameters. These initialization parameters are chosen by evaluating the merit function value for 10000 sets of five independently generated random numbers. The five random numbers
Figure 3. **Optimal solution:** The figure shows the target spectra and the obtained spectra of single layer disordered metamaterial found using the hybrid optimization technique. The host medium was found to be TiN with HfO2 particle inclusions having a radius of $\sim 34.73 \text{ nm}$ and a fill fraction $\rho \sim 0.84$. Subplots (a), (b) and (c), show the reflection, transmission and absorption spectra respectively.

The hybrid optimization algorithm is used on realistic target spectra. Reflection and transmission spectral features of a 5 layered multilayer stack consisting of arbitrary materials and thickness serves as realistic target spectra. The target spectra are shown in Fig. 3 in black. The above introduced hybrid optimization approach is utilized to find the optimal parameters of the single disordered metamaterial layer. Fig. 3 shows the optimal solution obtained utilizing this
The effective single layer permittivity and tuning ENZ: The subplot (a) shows the effective permittivity, $\varepsilon = \varepsilon' + \varepsilon''$ of the single disordered nanocomposite layer. The subplot (b) shows real part of permittivity $\varepsilon'$ calculated for various fill fractions, $\rho$. The inset figure shows the red shift in the ENZ position with increasing fill fraction, $\rho = 0$ and $\rho = 1$ show the bounds within which the ENZ can be varied.

The equivalent single layer solution consists of TiN host and $HfO_2$ particle inclusions. The effective thickness of this single layer is $\sim 97\, \text{nm}$ and host particles have a radius of $\sim 34.73\, \text{nm}$. We note here that this approach is not limited to the infrared spectral region but can also be utilized in the visible wavelength range. The effective layer’s permittivity is shown in Fig.4(a). The versatility of the proposed hybrid optimization is attributed to the ability of tuning the material dispersion properties of the disordered nanocomposite material over a range of wavelengths. The spectral position of epsilon near zero (ENZ) of the nanocomposite layer can be tuned by varying the fill fraction ($\rho$) of the inclusions. This is illustrated in Fig.4(b).

We perform abinitio FDTD simulations (Lumerical FDTD solutions) to find the reflection and transmission spectra of the actual disordered metamaterial. A spatial resolution of 1 nm and plane-wave source was used as an excitation source. We found no dependence on the polarization of incident light on the transmission and reflection spectra. Fig.5 shows the obtained transmission and reflection spectra. There is a qualitative agreement between EMT and FDTD. Low transmission in FDTD results compared to EMT is attributed to the fact that in EMT we assume a uniform, clear medium which is not the case with finite size inclusions. Furthermore, as the higher fill fraction values are not practically possible to achieve (without the spherical inclusions merging into each other), there is an upper bound on the practically achievable fill fraction. This leads to the spectral mismatch between EMT and FDTD results. However, there is a qualitative agreement between the two.

4. Conclusions

We have introduced the concept of spectral domain inverse design. This is complementary to existing techniques of structural optimization since we search for a complex dielectric permittivity. We have shown an approach to engineer new metamaterials that can be utilized in inverse design problems. Using this approach we have found a new disordered metamaterial that can be utilized as a TPV thermal emitter. Our work can lead to a new approach for metamaterials discovery.

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Figure 5. **FDTD simulations** The subplot (a) shows the transmission spectrum obtained using FDTD simulations. The subplot (b) shows the reflection spectrum obtained using FDTD simulations. This shows qualitative agreement with the EMT.

**Disclosures**

The authors declare no conflicts of interest.

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