A study of diffuse radio sources and X-ray emission in six massive clusters

V. Parekh,1* K. S. Dwarakanath,1 R. Kale2 and H. Intema3

1Raman Research Institute, C. V. Raman Avenue, Sadashivnagar, Bangalore 560080, India
2National Centre for Radio Astrophysics, T. I. F. R., Post Bag 3, Ganeshkhind, Pune 411007, India
3Leiden Observatory, Leiden University, PO Box 9513, NL-2300 RA Leiden, the Netherlands

ABSTRACT

The goal of this study is to extend our current knowledge of the diffuse radio source (halo and relic) populations to $z > 0.3$. Here, we report GMRT and EVLA radio observations of six galaxy clusters taken from the MAssive Cluster Survey (MACS) catalogue to detect diffuse radio emission. We used archival GMRT (150, 235, and 610 MHz) and EVLA ($L$ band) data and made images at multiple radio frequencies of the following six clusters – MACSJ0417.5–1154, MACSJ1131.8–1955, MACSJ0308.9+2645, MACSJ2243.3–0935, MACSJ2228.5+2036, and MACSJ0358.8–2955. We detect diffuse radio emission (halo or relic, or both) in the first four clusters. In the last two clusters, we do not detect any diffuse radio emission but we put stringent upper limits on their radio powers. We also use archival Chandra X-ray data to carry out morphology and substructure analysis of these clusters. We find that based on X-ray data, these MACS clusters are non-relaxed and show substructures in their temperature distribution. The radio powers of the first four MACS clusters are consistent with their expected values in the $L_X–P_{1.4\text{GHz}}$ plot. However, we found ultrasteepest spectrum radio halo in the MACSJ0417.5–1154 cluster whose rest-frame cut-off frequency is at $\sim 900$ MHz. The remaining two clusters whose radio powers are $\sim 11$ times below the expected values are most likely to be in the ‘off-state’ as has been postulated in some of the models of radio halo formation.

Key words: radio continuum: general – X-rays: galaxies: clusters.

1 INTRODUCTION

Giant radio haloes and relics are large ($\sim$Mpc size) diffuse synchrotron sources observed in massive and merging galaxy clusters. These extended radio sources are associated with diffuse, non-thermal, and low surface brightness radio emission with steep spectra ($<0.7 \gtrsim \alpha \gtrsim -1.4$; $S \propto v^{-\alpha}$, where $\alpha$ is the spectral index and $S$ is the flux density at frequency $v$). Radio haloes are generally centrally located in galaxy clusters and coincident with X-ray emission while radio relics are found at the peripheries of galaxy clusters (Feretti et al. 2012, and reference therein).

The formation and evolution of diffuse radio sources are still not clear. Several theories have been made for the mechanism of transferring energy into the relativistic electron population to generate cluster-wide radio haloes. The relativistic particles could be injected into the cluster volume by active galactic nucleus activity such as quasars and radio galaxies. This primary electrons’ population needs to be re-accelerated (Brunetti et al. 2001; Petrosian 2001) to compensate for radiative losses. A recent cluster merger is the most likely process to stimulate the re-acceleration of relativistic particles through injection of turbulence in the cluster volume (Brunetti et al. 2001; Petrosian 2001; Petrosian & Bykov 2008). Another model for radiating particles in radio haloes incorporates secondary electrons as the result of inelastic nuclear collisions which have occurred between the relativistic protons and thermal ions of the ambient ICM. The protons diffuse on a large scale because their energy losses are negligible. They can continuously produce in situ electrons which are distributed throughout the cluster volume (Blasi & Colafrancesco 1999; Miniati et al. 2001). It is possible that the high-energy electrons, responsible for the synchrotron emission, arise from the decay of secondary products of the neutralino annihilation in the dark matter haloes of galaxy clusters, and this has also been proposed (Colafrancesco & Mele 2001).

Different models have been suggested for the origin of the relativistic electrons radiating from relics. There is observational evidence that relics can be tracers of shock waves in merger events (Ensslin et al. 1998; Bagchi et al. 2006). These shocks, expanding with high velocity (Mach number $\sim 1–3$), can accelerate electrons to high energies and compress magnetic fields, giving rise to synchrotron radiation being emitted from large regions (Feretti et al. 2012). The accelerated particles will have a power-law energy distribution, as well as magnetic fields aligned...
parallel to the shock front. This is in agreement with their elongated structure almost perpendicular to the merger axis (van Weeren et al. 2011a, b).

It has been shown that X-ray data are crucial in the study of non-thermal radio emission from galaxy clusters. A number of authors have shown the correlation between radio luminosity of a halo ($P_{1.4\text{GHz}}$) versus the host cluster X-ray luminosity ($L_X$), mass, and temperature (Liang et al. 2000; Govoni et al. 2001; Feretti 2003; Cassano, Brunetti & Setti 2006; Cassano et al. 2007; Brunetti et al. 2009; Giovannini et al. 2009; Cassano, Brunetti & Venturi 2011; Venturi 2011). These correlations suggest that the most powerful radio haloes are found in the biggest and hottest clusters with the greatest X-ray luminosity. Furthermore, these correlations show a close link between non-thermal and thermal galaxy cluster physics. Govoni et al. (2001) showed the point-to-point correlation between X-ray and radio surface brightness for six radio halo clusters. This suggests a correlation between radio halo and X-ray emission in galaxy clusters.

High-resolution X-ray data showed a relationship between non-thermal radio sources and X-ray cluster morphology (Buote 2001; Schuecker et al. 2001; Cassano et al. 2010; Parekh et al. 2015, hereafter P15). This indicates a connection between the cluster dynamical state and the occurrence of a radio halo. The X-ray surface brightness distribution of merging clusters typically exhibits non-concentric iso-intensity contours, small-scale substructures, multiple peaks, and/or poor optical and X-ray alignments. Buote (2001) first showed a correlation between the power ratio ($P_1/P_0$) calculated for ROSAT observed ($\sim 14$) X-ray clusters and $P_{1.4\text{GHz}}$ measured for radio data. The power ratios measure both the multipole expansion of a 2D structure and the dynamical state of a cluster (Buote & Tsai 1995, 1996). Buote concluded that approximately $P_{1.4\text{GHz}} \propto P_1/P_0$, which means the clusters that host the most powerful radio haloes are undergoing the greatest exodus from a virialized state. Recently, Cassano et al. (2010) used three parameters, namely centroid shift (Mohr, Fabricant & Geller 1993; Poole et al. 2006; Maughan et al. 2008; Böhringer et al. 2010), third-order power ratio ($P_3/P_0$) (Weißmann et al. 2012), and concentration (Santos et al. 2008) to demonstrate a relationship between cluster mergers and the presence of a radio halo. The analysis of Cassano et al. (2010) supports the relationship between cluster dynamical activity (or dynamical state) and cluster-wide diffuse sources. This relationship, in turn, supports the particle (re-)acceleration mechanism in the formation of radio haloes in galaxy clusters. The current scenario indicates that diffuse radio features (radio haloes) are seen only in merging clusters.

### 1.1 Massive Cluster Survey

The Massive Cluster Survey (MACS) cluster sample comprises a total of 124 clusters in the redshift range $0.3 < z < 0.7$ (Ebeling, Edge & Henry 2001). These clusters are spectroscopically confirmed and represent a statistically complete sample of X-ray luminous [$\sim 10^{45}$ erg s$^{-1}$ (0.1–2.4 keV)], massive ($10^{14}$–$10^{15}$ M$_\odot$) and distant clusters of galaxies. Many MACS clusters have been studied for diffuse radio sources, for example MACS J0717+3745, MACSJ1149.5+2223, MACSJ1752.1+4440, and MACS J0535.4+3342 (Bonafede et al. 2009, 2012; van Weeren et al. 2009). The detection probability of radio halo is rather small (< 10 per cent) for all clusters, and the probability increases to ~40 per cent for clusters with $L_X > 10^{27}$ erg s$^{-1}$ (Venturi et al. 2008). Under these circumstances, majority of MACS clusters might be expected to host detectable radio haloes or relics, or both. With these criteria of most massive and luminous clusters, we selected a sample of six most massive ($> 6 \times 10^{14} M_\odot$) and disturbed galaxy clusters from the MACS catalogue (Ebeling et al. 2010). The Sunyaev–Zel’dovich (SZ) mass of these clusters is also very high which suggests that they are good candidates to host diffuse radio sources (Cassano et al. 2013). In this work, we studied these clusters, in radio and X-ray observations, to detect diffuse radio sources and their connection with cluster merger. We list the sample of MACS clusters in Table 1 with other X-ray properties.

This paper is organized as follows. Section 2 gives a brief overview of radio observations and data reduction. Section 3 gives the details about individual clusters in the sample. In Section 4, we present X-ray data, temperature map, and morphology analysis. Section 5 gives substructure analysis of six radio halo clusters. Finally, Section 6 gives discussion and conclusions. We assumed $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_M = 0.3$, and $\Omega_{\Lambda} = 0.7$ throughout the paper.

### 2 RADIO OBSERVATIONS AND DATA REDUCTION

#### 2.1 610 and 235 MHz GMRT data

We used GMRT archival data to image six MACS clusters. These clusters were observed during 2010 October. Each of the MACS cluster was observed for a total duration of ~8 h in the dual frequency (235/610 MHz) mode. This allows us to record one polarization at each frequency band. The GMRT software backend with bandwidths of 6 and 32 MHz at 235 and 610 MHz

| MACS clusters | Alt. name | RA J2000 | Declination J2000 | z | Scale (kpc arcsec$^{-1}$) | $Y_{SZ,M_{500}}$ ($10^{14}$ M$_\odot$) | $T$ (keV) | $L_X$ (10$^{45}$ erg s$^{-1}$) |
|---------------|-----------|----------|-------------------|---|-----------------------------|---------------------------------|----------|-----------------------------|
| MACSJ0417.5−1154 | A1300 | 04 17 34.6 | −11 54 32 | 0.44 | 5.68 | 12.25 ±0.52 | 0.55 | 11.07 | 3.66 |
| MACSJ1131.8−1955 | A1300 | 11 32 00.7 | −19 53 34 | 0.30 | 4.45 | 8.97 ±0.45 | 0.45 | 7.75 | 1.54 |
| MACSJ0308.9+2645 | 03 08 55.9 | +26 45 38 | 0.35 | 4.94 | 10.75 ±0.62 | 0.65 | 10.54 | 1.79 |
| MACSJ2243.3−0935 | 22 43 21.1 | −09 35 43 | 0.44 | 5.68 | 10.00 ±0.43 | 0.44 | 7.98 | 1.89 |
| MACSJ2228.5+2036 | RXJ2228.6+2037 | 22 28 37.1 | +20 36 31 | 0.41 | 5.45 | 8.26 ±0.43 | 0.44 | 8.16 | 1.63 |
| MACSJ0358.8−2955 | A3192 | 03 58 42.5 | −29 54 30 | 0.42 | 5.53 | 7.20 ±0.52 | 0.49 | 8.8 | 2.35 |
frequencies, respectively, was used in this observation. Data were analysed using the Astronomical Image Processing System and Common Astronomy Software Applications (CASA) packages [developed by National Radio Astronomy Observatory (NRAO)]. The data were inspected for the radio frequency interference, non-working antennas, bad baselines, channels, and time period. Corrupted data were excised from the u-ν data set. The flux density of each primary or flux calibrator(s) is set according to the Baars et al. (1977). The same calibrator was used for the bandpass calibration followed by determining the flux density of the secondary or phase calibrator(s) using the antenna gain solutions. Calibrated visibilities then used to create the images using the standard Fourier transform deconvolution method. Few rounds of self-calibration (three phases + one amplitude) were applied to reduce the effects of residual phase errors in the data and to improve the quality of the final images. Wide-field imaging with 256 projection planes was used in the CASA task ‘clean’ (see CASA manual). Images were produced with a variety of weighting (uniform and natural) schemes for the visibilities.

2.2 150 MHz TGSS data

The TGSS is a fully observed yet largely unreleased survey of the radio sky at 150 MHz as visible from the GMRT, covering the full declination range of −55 to +90 deg. All these observational data were recorded in half polarization (RR, LL) every 2 s in 256 frequency channels across 16 MHz of bandwidth (140–156 MHz). Each pointing was observed for about 15 min, split over three or more scans spaced out in time to improve UV coverage. As a service to the community, these archival data have been processed with a fully automated pipeline based on the SPAM package (Intema et al. 2009; Intema 2014), which includes direction-dependent calibration, modelling, and imaging to suppress mainly ionospheric phase errors. In summary, the pipeline consists of two parts: a preprocessing part that converts the raw data from individual observing sessions into pre-calibrated visibility data sets for all observed pointings and a main pipeline part that converts pre-calibrated visibility data per pointing into Stokes I continuum images. The flux scale is set by calibration on 3C48, 3C147, and 3C286 using the models from Scaife & Heald (2012). More details on the processing pipeline and characteristics of the data products will be given in an upcoming paper on the first TGSS alternative data release (Intema et al. 2016; ADRI). For this study, ADRI images were used to create mosaics at the cluster positions. All the 150 MHz images, in our study, have a resolution of 20 arcsec × 20 arcsec.

2.3 L-band ELVA data

We used EVLA archival observations for the sample of MACS clusters. In the archive, we found L band (1400 MHz), C configuration (max baseline: 3.4 km) observations 11B-018 (2012 February) for only two clusters: MACSJ0417.5−1154 and MACSJ2243.3−0935. Each of these clusters was observed for a duration of ~1.5 h. These observations were carried out with 16 spectral windows (spws) from 1194 to 1892 MHz. Each spectral window has 64 channels. About half of the data were corrupted, so we removed them from the u-ν data. We followed the standard procedure of data analysis. The flux density of each primary or flux calibrator(s) is set according to the Perley & Butler (2013). The calibrated visibilities were then imaged in CASA in the same way as described for 235 and 610 MHz observations. In further analysis, we convolved all images (235, 610, and 1575 MHz) to 20 arcsec × 20 arcsec resolution, which is the resolution of the 150 MHz TGSS images.

In order to make an error estimate in the flux density measurements, here we outline the procedure. There are two primary sources of errors in the flux density measurements: (1) an error due to the uncertainties in the flux densities of the unresolved source(s) used for calibration of the data – we assumed this error to be ~10 per cent – and (2) since diffuse radio sources are extended sources, the errors in their flux density estimations will be the rms in the image multiplied by the square root of the ratio of the solid angle of the source to that of the synthesized beam. These two sources of errors are unrelated, so they are added in quadrature to estimate the final error on the flux densities of the extended sources, as given below:

\[ \Delta S = \left( \sigma_{\text{amp}} S^2 + \sigma_{\text{rms}} \sqrt{n_{\text{beams}}} \right)^{1/2}, \]

where \( \Delta S \) is a flux density, \( \sigma_{\text{amp}} \) is the image rms noise, and \( n_{\text{beams}} \) is the number of beams in the extent of the source.

3 INDIVIDUAL CLUSTER

3.1 MACSJ0417.5−1154

MACSJ0417.5−1154 is a hot (~11 keV), X-ray luminous (3.66 × 10^{45} \text{ erg s}^{-1}), and massive cluster (12.25 × 10^{14} \text{ M}_\odot) at a redshift of \( z = 0.44 \). MAC0417.5−1155 is the most massive cluster in the sample. The detection of a halo in this cluster was first reported by Dwarakanath, Malu & Kale (2011). Fig. 1 shows the (a) 235 MHz, (b) 610 MHz GMRT radio images, and (c) 1575 MHz EVLA radio contours overlaid on the Chandra X-ray image of MACSJ0417.5−1155. In all these frequencies, there is extended, comet-like diffuse radio emission visible along the north-west direction. As shown in Fig. 1(c), diffuse radio emission is coincident with the corresponding hot X-ray distribution in direction and size. Elliptical X-ray morphology clearly suggests that it is a disturbed and merging cluster. In 610 MHz observation, the largest linear size of the halo is ~1 Mpc which is the typical size of radio haloes in massive clusters. In the X-ray image, a bright and compact core is visible which is situated with the optical core of one of the two merging clusters (Mann & Ebeling 2012, hereafter AM12). Furthermore, this morphology is also similar to those observed in merging clusters such as A2146 and Bullet Clusters, which indicates a high-velocity merger (AM12).

Table 2 contains flux densities and linear sizes of haloes (within 3σ contours) measured at 235, 610, and 1575 MHz. There are two unresolved point sources (A and B) embedded in the diffuse radio emission (Fig. 1). To calculate the integrated flux density of the radio halo, we subtracted flux densities of these point sources (Table 3), which are calculated by Gaussian fitting, from the total halo flux density measured within the given size at each frequency in Table 2. Fig. 1(d) shows the spectrum of the radio halo of MACSJ0417.5−1155. Between 235 and 610 MHz, the spectrum is flat (\( \alpha_1 = -0.38 \)), and between 610 and 1575 MHz the spectrum is steep (\( \alpha_2 = -1.72 \)). We also compared EVLA 1575 MHz flux density measurements with the L-band NVSS (Condon et al. 1998) observation of MACSJ0417.5−1154. In the NVSS image, A and B point sources blended, so we calculated flux density for the whole halo region and then subtracted flux densities of A and B derived from the EVLA observation with an assuming spectral index of 0.7.

1. TIFR GMRT Sky Survey; see http://tgss.ncra.tifr.res.in.

2. In this work, we scaled all values of flux densities of haloes and relics to the Baars scale according to Baars et al. (1977).
Figure 1. MACSJ0417−1155. (a) GMRT 235 MHz contours, (b) GMRT 610 MHz contours, (c) EVLA 1575 MHz contours on Chandra X-ray image and (d) the integrated spectrum of the radio halo in MACSJ0417−1155. All radio images have a resolution of 20 arcsec × 20 arcsec. First contour is drawn at 3σ where σ_{235 MHz} = 0.44 mJy beam$^{-1}$, σ_{610 MHz} = 0.15 mJy beam$^{-1}$, and σ_{1575 MHz} = 40 µJy beam$^{-1}$. Contour level increases in steps of $\sqrt{2}$. Dashed line shows −3σ contours. The total flux density of the halo is $\sim$18 mJy, which is comparable with the EVLA result. Spectral index images of radio haloes are very important to understand the presence of any magnetic field gradients and/or spatial variations of the relativistic electrons high-energy cut-off. But, signal-to-noise ratios in these radio images are not adequate to produce a spectral index map. There is no diffuse emission visible in the 150 MHz TGSS data. The expected surface brightness ($\sim$2 mJy beam$^{-1}$) is below the rms ($\sim$4 mJy beam$^{-1}$) of 150 MHz observation.

We estimate the equipartition magnetic field (see Carilli & Taylor 2002; Govoni & Feretti 2004; Ferrari et al. 2008, for reviews) in the MACSJ0417−1155 radio halo to be $\sim$1.0 µG.

3.2 MACSJ1131.8−1955

MACSJ1131.8−1955 or A1300 (z = 0.30) is a complex and post-merging cluster as evident by both X-ray and optical observations (Lemonon et al. 1997; Ziparo et al. 2012). Its X-ray luminosity is $1.54 \times 10^{45}$ erg s$^{-1}$ and mass is $7.75 \times 10^{14}$ M$_{\odot}$. First detection of the radio halo and relic in MACSJ1131.8−1955 was reported by Reid et al. (1999). They have detected diffuse radio sources at higher frequencies (1.34, 2.4, 4.8, and 8.6 GHz) using ATCA and MOST radio telescopes. It was also detected in the 325 MHz GMRT observations (Venturi et al. 2013). Here we report halo and relic detection at 150, 235, and 610 MHz observations using the GMRT.

Fig. 2 shows (a) 150 MHz TGSS, (b) 235 MHz GMRT, and (c) 610 MHz GMRT observations superimposed on the Chandra X-ray image of MACSJ1131.8−1955. In all three frequencies, radio halo and south-west relic are detected. This south-west relic is marked with a ‘+’ sign in Fig. 2(c). As we noticed, in all three bands, radio halo extends mainly in the east direction of the peak in the X-ray surface brightness distribution. Furthermore, radio halo looks more extended (in the east direction) in 610 MHz data as compared to the 150 and 235 MHz observations. It has very disturbed and elongated X-ray morphology which indicates a series of current mergers. The overall shape and size of both halo and relic are comparable to the latest findings in 325 MHz observation. As we detected, largest linear sizes of halo and relic at 610 MHz are $\sim$815 and $\sim$446 kpc, respectively, which are comparable to the largest linear sizes of halo $\sim$890 kpc and relic $\sim$450 kpc at 325 MHz, respectively. In 150 MHz observation, we detected an elongated structure (above 3σ contour)
at the position of south-west relic as seen in 235 and 610 MHz observations. We marked two unresolved point sources (A and B) in the 150 MHz image which are also visible at the same position in the 235 and 610 MHz images. Since diffuse radio emission is very complex, it is difficult to avoid these point sources to calculate total flux density of the radio halo. The integrated flux densities (after subtracting flux densities of point sources) of the radio halo and relic, within the $3\sigma$ contour level, are reported in Table 2.

We show halo and relics spectra in Fig. 3. We used NVSS flux density (20 mJy) of the radio halo and relic. The integrated spectral index for the halo is $\alpha_{1400\text{MHz}} = 1.37 \pm 0.11$ which is comparable with finding of Giacintucci (2011). If we use ATCA flux density (10 mJy) then $\alpha_{1400\text{MHz}} = 1.72 \pm 0.12$. The spectral index for the south-west relic is $\alpha_{1400\text{MHz}} = -1.06 \pm 0.13$ which is comparable with the previous findings (Giacintucci 2011; Venturi et al. 2013).

Table 3. Flux densities of unresolved sources in MACS clusters.

| Cluster | Frequency (MHz) | Flux density of unresolved point sources (mJy) |
|---------|----------------|-----------------------------------------------|
| MACSJ0417.5−1154 | 235 | 46 ± 4.6 17 ± 1.7 (A and B) |
| | 610 | 30 ± 3.6 ± 0.6 |
| | 1575 | 13 ± 1.34 ± 0.4 |
| MACSJ1131.8−1955 (halo) | 150 | 357 ± 37 114 ± 11 (A and B) |
| | 235 | 580 ± 58 190 ± 19 |
| | 610 | 85 ± 8.5 ± 3 |
| MACSJ0308.9+2645 | 610 | 30 ± 3 42 ± 4.2 3 ± 0.3 (A, B, and C) |

Venturi et al. (2013) have noticed candidate relic ($\sim 700$ kpc size) in the 325 MHz image, based on its overall morphology and location, at the north-west direction from the centre of cluster. We also noticed this candidate relic in 610 MHz observation, which is also marked with a ‘+’ sign. We measured the total flux density of this candidate relic to be $11 \pm 1.4$ mJy and the largest linear size is $\sim 670$ kpc which are comparable with the previous results. We show its spectrum between 325 and 610 MHz in Fig. 3. The spectrum ($\alpha = -1.17$) of the north-west relic is steeper than that of the south-west relic. This candidate relic is not visible in the 150 and 235 MHz images. The sensitivities of latter observations are not adequate to detect this candidate relic. In Fig. 2(c), we notice a bridge of radio emission between the central halo and the south-west relic in the 610 MHz image. This is one of the rare examples where a radio bridge is clearly visible between the halo and the relic. Due to lack of sensitivity, this bridge is not visible in the 150 and 235 MHz observations.

We estimate the equipartition magnetic fields in the MACSJ1131.8−1955 radio halo and (south-west) relic to be $\sim 1$ and $\sim 1.3 \mu$G, respectively.

3.3 MACSJ0308.9+2645

MACSJ0308.9+2645 is a massive ($10.75 \times 10^{14}$ M$_{\odot}$), luminous (1.79 $\times 10^{45}$ erg s$^{-1}$), hot (10.54 keV) cluster situated at $z = 0.35$. We analysed 150, 235, and 610 MHz GMRT data to search for diffuse radio emission in this cluster. We detected candidate radio halo, only in 610 MHz observation. Fig. 4 shows the (a) 610 MHz image and (b) 610 MHz contours overlaid on the Chandra X-ray image. Its X-ray morphology seems to be relaxed and regular. However, its temperature distribution shows two arc-like substructures (see Section 5), and radio halo is situated between these substructures. The largest linear size of the candidate radio halo is $\sim 850$ kpc. There are three point sources (A, B, and C) embedded in the
Figure 2. MACSJ1131.8−1955. (a) TGSS 150 MHz contours, (b) GMRT 235 MHz contours, and (c) GMRT 610 MHz contours on Chandra X-ray image. All radio images have a resolution of 20 arcsec × 20 arcsec. First contour is drawn at $3\sigma$ where $\sigma_{150\text{MHz}} = 5 \text{ mJy beam}^{-1}$, $\sigma_{235\text{MHz}} = 3 \text{ mJy beam}^{-1}$, and $\sigma_{610\text{MHz}} = 0.30 \text{ mJy beam}^{-1}$. Contour level increases in steps of $\sqrt{2}$. The + sign is for the relics sources. The dashed line shows $-3\sigma$ contours.

Figure 3. The integrated spectrum of halo (left) and relic(s) (right) of MACSJ1131.8−1955.
diffuse radio source. We report 610 MHz integrated flux density (after subtracting flux densities of point sources) of the halo within 3σ contour level in Table 2. There is no diffuse radio emission visible in the 150 and 235 MHz images. The expected surface brightness of the radio halo is \( \sim 3.3 \) mJy beam\(^{-1}\) at 150 MHz, and \( \sim 1.8 \) mJy beam\(^{-1}\) at 235 MHz. Both these expected values are below the rms value (\( \sim 4 \) mJy beam\(^{-1}\)) in these frequencies. We estimated the radio power at 1400 MHz (\( P_{14\mathrm{GHz}} \)) by extrapolating the flux density at 610 MHz to the flux density at 1400 MHz using a spectral index of \( -1.3 \) (which is the average value of the spectral index of the sample of haloes given by Feretti et al. 2012). The equipartition magnetic field in the MACSJ0308.9+2645 radio halo is \( \sim 0.7 \) \( \mu \)G.

### 3.4 MACSJ2224.3−0935

This cluster is located at the centre of a supercluster of SCL2224.3−0935, at a redshift of \( z = 0.44 \). Three massive filaments of size (5–15) \( h_0^{-1} \) Mpc and mass of 1.53 \( \pm 1.01 \times 10^{15} \) M\(_\odot\) have been discovered around this cluster. The size of this supercluster is estimated to be \( 45 \times 15 \times 30 \) \( h_0^{-1} \) Mpc (\( l \times w \times d \)), making it one of the largest superclusters known as intermediate redshifts (Schirmer et al. 2011). There are \( \sim 14 \) cluster members embedded in this large filamentary cosmic web network. As compared to other superclusters such as A901/902, MS0302, or CL0016, SCL2224.3−0935 is significantly larger in size, mass, and host member clusters. The core cluster of SCL2224.3−0935, MACSJ2224.3−0935, is massive (\( M_{200} = 1.54 \pm 0.29 \times 10^{15} \) M\(_\odot\)) and shows significant substructuring. It contains several strongly lensed galaxies and is not in dynamical equilibrium based on galaxies distribution (Schirmer et al. 2011).

The detection of a radio halo in this cluster was first reported by Cantwell et al. (2016). We analysed its 235 and 610 MHz GMRT data, and 1400 MHz EVLA data to study the diffuse radio emission. No diffuse radio emission is visible in the 150 (TGSS), 235, and 1400 MHz data. Fig. 5 shows (a) the 610 MHz image for the larger region around the core of the cluster and (b) 610 MHz contours on the Chandra X-ray image for the central core region of the cluster. We detected a candidate radio halo of largest linear size of \( \sim 722 \) kpc, and perpendicular to the merger axis in the 610 MHz image. We measured the total flux density of the radio halo to be \( 9 \pm 1.03 \) which is comparable with the findings of Cantwell et al. (2016).

X-ray observations clearly show a disturb morphology. There are many active radio galaxies visible around the cluster, mainly in the east direction from the cluster centre. Some of them have one-sided or two-sided jet emission. Some of them have optical counterparts in the SDSS survey. Cantwell et al. (2016) have also detected a candidate relic in MACSJ2243.3−0935 using GMRT 610 MHz newer observation. This candidate relic is marked with a ‘+’ sign in Fig. 5. This relic is situated in the west direction at a distance of \( 2.13 \) Mpc from the cluster centre. We estimated its largest linear size, \( \sim 680 \) kpc, which is the same as Cantwell et al. (2016). However, our flux density measurements of 2.0 \( \pm 0.4 \) mJy are not consistent with their measurements. More details on discrete radio sources around the radio halo are given in Cantwell et al. (2016). In the X-ray image (Fig. 5b), there are two subclusters in the west and east directions, and diffuse radio emission is situated in between them. As shown in fig. 4 of Schirmer et al. (2011), MACSJ2224.3−0935 has two clumps of galaxy distribution. This suggests that there are two clusters which are merging or have already passed through each other (post-merger stage).

We report 610 MHz integrated flux density of the halo in MACSJ2224.3−0935 measured within 3σ contour level in Table 2. We estimated expected surface brightness values for the halo at 150 (~2.11 mJy beam\(^{-1}\)), 230 (~1.18 mJy beam\(^{-1}\)), and 1400 MHz (~0.11 mJy beam\(^{-1}\)). These expected values are at the level of 2–3\( \sigma \) in these observations. We estimated the radio power at 1400 MHz (\( P_{14\mathrm{GHz}} \)) by extrapolating the flux density at 610 MHz to the flux density at 1400 MHz using a spectral index of \( -1.3 \). The equipartition magnetic field in this candidate radio halo is \( \sim 0.7 \) \( \mu \)G.

### 3.5 MACSJ2228.5+2036 and MACSJ0358.8−2955

These two MACS clusters are massive and hot, as well as disturbed in their X-ray morphology, but there is no diffuse radio emission visible at any of the observed frequencies (Table 2). Venturi et al. (2008) have first observed MACSJ2228.5+2036 using GMRT as a...
part of GMRT radio halo survey. They have also not detected any diffuse radio emission in it. Hence, we calculate the upper limit on the radio luminosity of the haloes that could be present in these two clusters. We report upper limits for these two clusters in Table 2. We calculated the upper limit (at 610 MHz) by first measuring the rms in the central regions of these two clusters. This rms is measured for a single beam, so we need to scale it with a factor of $\sqrt{2}$ to take account of the number of beams within the halo size. We used a canonical size for the radio haloes to be $\sim 1$ Mpc. An upper limit to the radio power of a halo at 1400 MHz can be estimated using the upper limit at 610 MHz. We used a spectral index of $-1.3$ to extrapolate the upper limit at 610–1400 MHz.

4 X-RAY DATA ANALYSIS AND IMAGES

In this work, we used Chandra archival data for the sample of six MACS clusters (Table 4). We processed these data with CIAO 4.6 and CALDB 4.6.1.1. We first used the chandra_repro task to reprocess all ACIS imaging data, removed any high background flares (3$\sigma$ clipping) with the task lc_sigma_clip, followed by combining multiple data sets using merge_obs script. All event files included the 0.3–7 keV broad energy band and 2 arcsec pixels binning. We removed point sources around the cluster. We divided counts’ image with the exposure map and generated the flux image (photon cm$^{-2}$ s$^{-1}$). Finally, we smoothed this image with $\sigma = 10$ arcsec to remove zero counts.

4.1 Temperature maps

To understand the hot ICM morphology and its relation with diffuse radio sources, we generated temperature maps of MACS clusters. The temperature map of the ICM is also a good indicator of the dynamical activity in the cluster. We do not show temperature maps for the non-detection radio halo clusters, i.e. MACSJ2228.5+2036 and MACSJ0358.8−2955. For the remaining clusters, we used the highest exposure time data. We used X-ray Monte Carlo (XMC) technique to generate the temperature map (see Andersson, Peterson & Madejski 2007; Peterson, Marshall & Andersson 2007; Andersson et al. 2009, for details of the method). In this technique, it is assumed that the given intracluster hot gas is composed of a number of Gaussian ‘blobs’ or smooth particles. Each blob has its own temperature, flux, abundance, electron density, phi and psi (sky location), Gaussian width (size), redshift, absorption column, etc. X-ray emission from these hypothetical ‘blobs’ is then propagated through the (simulated) telescope’s instrument, generating dummy X-ray events in a similar way to that which would come from an actual X-ray observation. These simulated events are compared with the real data events being modelled. In this method, the spectral and spatial models are used together. We used the warm-absorbed APEC (spectral) model along with the Chandra detector (spatial) model derived for the XMC. In this procedure, given (free) parameters (APEC model parameters such as temperature and normalization, and spatial coordinates of detectors) of the ICM are iterated using the Markov Chain Monte Carlo (MCMC) technique. Table 5 shows the parameters and their ranges. For instance, in this analysis we allowed the temperature to vary between $1$ and 15 keV, while fixed the solar abundance to 0.3 for all clusters. We used the NRAO H1 all sky survey data (Dickey & Lockman 1990) to derive the $n_{H}$ values for each cluster.

When all parameters converge after a sufficient number of iterations, the resulting distribution of blobs describes the shape and characteristics of the galaxy cluster under investigation (Andersson et al. 2007). The final results have a number of statistical samples of acceptable or converged fits which fall into the ‘confidence region’ of the ICM’s input parameters. These ‘well-fitted’ parameters

![Figure 5. MACSJ2243.3−0935. (a) GMRT 610 MHz contours (a resolution of 20 arcsec × 20 arcsec) for a larger region around the core of the cluster and (b) GMRT 610 MHz contours on Chandra X-ray image for the central halo region. First contour is drawn at 3$\sigma$ where $\sigma_{610\text{MHz}} = 0.10$ mJy beam$^{-1}$. Contour level increases in steps of $\sqrt{2}$. The + sign is for the candidate relic. The dashed line shows $-3\sigma$ contours.](https://academic.oup.com/mnras/article-abstract/464/3/2752/2447925)
Table 5. Spectral and spatial model parameters and their values for XMC analysis.

| Parameters       | Fixed/free | Min and max values |
|------------------|------------|--------------------|
| Spectral model   |            |                    |
| $n_H$            | Fixed      | NRAO H1 all sky survey |
| $T$ (keV)        | Free       | 1–15 keV           |
| $Z$⊙             | Fixed      | 0.3                |
| $z$              | Fixed      | According to Table 1|
| Spatial model    |            |                    |
| $x$              | Free       | −8.4 to 8.4 arcmin |
| $y$              | Free       | −8.4 to 8.4 arcmin |
| ln $\sigma$ (arcsec) | Free    | 0–6 arcsec          |

describe the properties of the ICM. In this work, all the results derived from the model samples are from the iterations of 500–3000, where the value of $\chi^2$ is reduced (to $\sim$1) and stable, and considered to be a converged chain where all corresponding parameters have the best-fitting value. Figs 6 and 7 show the temperature maps of four MACS clusters. Temperature maps show high values because in the XMC method multitemperature plasma overlaps on each other which mimics the ‘real’ situation in ICM plasma (Andersson et al. 2009). The advantage of the XMC technique is that it explores, with the MCMC method, the extremely high dimensionality of parameter space of the blobs for a given ICM model. The MCMC technique efficiently samples the probability distribution of a model’s parameters without being trapped in local likelihood maxima. The result is a well-fitted list of sample cluster models, all of which are consistent with the data.

4.2 X-ray morphology analysis

Substructures, X-ray morphology, and/or X-ray centroid variations are typical features that allow us to understand whether or not a cluster is virialized. The X-ray-emitting gas in galaxy clusters carries signatures of dynamical activity that manifests itself as distortions in the X-ray surface brightness images. We computed three non-parametric morphology parameters, $Gini$, $M_{20}$, and concentration ($C$) using the Chandra X-ray images of MACS clusters to characterize the degree of disturbances in its ICM. These three parameters are known to be effective in segregating galaxy clusters according to the level of disturbances in them (P15). $Gini$ is a measure of the flux distribution among the image pixels; its value is 0 if the flux is equally distributed among the pixels (typically non-relaxed clusters) and $\sim$1 if most of the flux is contained only in a small number of pixels (typically relaxed and cool-core clusters; Lotz et al. 2004).

The moment of light, $M_{20}$, is the normalized second-order moment of the relative contribution of the brightest 20 per cent of the pixels (Lotz et al. 2004) and is a measure of the spatial distribution of bright cores and substructures in the cluster. The typical values of $M_{20}$ are between $-2.5$ (very relaxed) and $-0.7$ (very disturbed) (see P15). The parameter $C$ is a measure of the concentration of the flux in the cluster that depends on the ratio of the radii at which 80 and 20 per cent of the cluster flux are found (Conselice 2003) and has a minimum value of 0.0 which indicates the most disturbed clusters.

For these MACS clusters we used the 500 kpc region around the cluster centroid to calculate the morphology parameters. We have listed the morphology parameters in Table 6 along with $1\sigma$ uncertainties. We refer the reader to P15 and the references therein for the details of the calculations of X-ray morphology parameters. We have plotted each of these parameters versus other two in Fig. 8. In
MACS clusters

Figure 7. (Left) MACSJ0308.9+2645 and (right) MACSJ2243.3−0935 (for only central halo region) temperature maps. In both these images, contours are corresponding to 610 MHz GMRT radio observations. Contour levels are the same as in Figs 4 and 5. Colourbar is in units of keV. We displayed temperature values at ∼30 per cent confidence level.

Table 6. MACS clusters morphology parameter values.

| Cluster       | Gini     | $M_{20}$ | Concentration |
|---------------|----------|----------|---------------|
| MACSJ0417.5−1154 | 0.55 ± 0.01 | −1.72 ± 0.34 | 1.23 ± 0.18   |
| MACSJ1131.8−1955  | 0.46 ± 0.00  | −1.52 ± 0.30  | 0.07 ± 0.37   |
| MACSJ0308.9+2645  | 0.49 ± 0.00  | −1.77 ± 0.29  | 0.29 ± 0.18   |
| MACSJ2243.3−0935  | 0.30 ± 0.00  | −1.14 ± 0.22  | 0.74 ± 0.39   |
| MACSJ2228.5+2036  | 0.47 ± 0.00  | −1.62 ± 0.37  | 0.13 ± 0.56   |
| MACSJ0358.8−2955  | 0.62 ± 0.00  | −1.53 ± 0.34  | 1.39 ± 0.62   |

5 SUBSTRUCTURE ANALYSIS

As discussed in Section 4.2, all the MACS clusters in this sample appear to be dynamically disturbed. The morphological parameters of these galaxy clusters quantify the degree of disturbances in the ICMs of the respective clusters. We compare the parameters estimated for MACS clusters with those for a sample of 85 clusters with deep Chandra observations for which morphological parameters are available in P15. Among the sample of 85 clusters, there are 49 low-z (0.2–0.3) and 36 high-z (0.3–0.8) clusters with X-ray flux, $f > 1.4 \times 10^{-13} \text{ erg s}^{-1} \text{ cm}^{-2}$ and $L_X > 3 \times 10^{43} \text{ erg s}^{-1}$. P15 have classified the clusters, based on the combination of morphological parameters, into four categories of their dynamical states:

- Strong relaxed (SR): $C > 1.55$, $M_{20} < −2$, and $Gini > 0.65$.
- Intermediate clusters: relaxed (R) and non-relaxed (NR): $1.0 < C < 1.55$, $−2.0 < M_{20} < −1.4$, and $0.40 < Gini < 0.65$.
- Strong non-relaxed (SNR): $C < 1.0$, $M_{20} > −1.4$, and $Gini < 0.40$.

In Fig. 8, the points corresponding to the six MACS clusters are plotted along with those for the P15 sample of clusters (Giovannini et al. 2009) in the morphological parameter planes. The SR clusters are well separated from SNR clusters, while intermediate clusters are found in between them. The morphological parameters (P15) for the sample of radio halo clusters from Giovannini et al. (2009) are also shown in Fig. 8 (hexagons). The clusters with radio haloes are mainly in the NR (67 per cent) and SNR (33 per cent) categories, while the MACS clusters are in the intermediate category. Furthermore, except the MACSJ2243.3−0935 (SNR), other five MACS clusters are NR. Fig. 9 shows the distribution of these clusters in the temperature and morphology parameter planes. Note that all the MACS clusters are in region 3 which mostly contains clusters with radio haloes.

We have also analysed the temperature maps of MACS clusters. It is believed that the distribution of temperature provides crucial information about the underlying substructure in the ICM which is sometimes masked in the surface brightness images. Furthermore, the reason for using the temperature maps is that, we see hot regions caused by a currently propagating or just-passed...
merger shock (Markevitch et al. 2003, 2005). In these shock-heated regions, gas expands adiabatically, reaching pressure equilibrium with the surrounding environment after the passage of the shock. We also see patchy temperature distribution, spiral, and whirlpool structures because of large-scale gas motions during the mergers causing subsequent mixing of gases at different temperatures. This re-distribution of temperature depends on whether the merger has disrupted the structure of the cool core, and the time elapsed since the merger. We have analysed the temperature maps of the MACS clusters in our sample along with their radio counterparts. Below we will discuss the substructure analysis of the sample of six MACS clusters.

MACSJ0417.5−1154. AM12 classified MACSJ0417.5−1154 as a binary head-on merger of similar mass. Temperature map [Fig. 6 (left)] shows disturbances largely in the outer parts of the cluster, while the inner core (∼100 kpc) does not appear to be affected by the merger. The morphology of the temperature distribution is similar to that of the diffuse radio emission.

MACSJ1131.8−1955. Its complex X-ray morphology indicates a series of recent mergers, as well as an ongoing merger (Lemonon et al. 1997; Reid et al. 1999; Ziparo et al. 2012). AM12 classified this cluster as the most complex mergers that have recently passed, or are still undergoing multiple mergers along various axes. Temperature map of MACSJ1131.8−1955 [Fig. 6 (right)] is irregular and shows a complex substructure in the core. In the temperature map, we observed large-scale gas motions as similar to the whirlpool structure noticed by Finoguenov, Böhringer & Zhang (2005). Hot gas (∼10 keV) is extended in the north-east direction where most of the radio emission is located. Recently, Ziparo et al. (2012) noticed a possible shock front (with a Mach number of $M = 1.20 \pm 0.10$) in the southern part of the cluster where the south-west relic is located. We also noticed, in the temperature map, a sharp edge near the position of the south-west relic. There is no evidence for a shock front near the candidate relic towards the north-west.

MACSJ0308.9+2645. Its X-ray surface brightness distribution is regular and smooth; however, its morphology parameters indicate that it is an NR cluster. The temperature map of MACSJ0308.9+2645 [Fig. 7 (left)] is rather complex in conformity with it being an NR cluster. These substructures are not visible in the corresponding surface brightness image (Fig. 4). There are two cool cores visible in the temperature map. It is possible that these two cores are part of merging clusters. Diffuse radio emission (at 610 MHz) is visible in between these two cool cores. Radio emission is extending in the north–south direction which is roughly perpendicular to the merger axis along the direction joining the cool cores. There is some extension of radio emission towards the east direction. Overall morphology of MACSJ0308.9+2645 indicates that it is an ongoing merger.

MACSJ2243.3−0935. It is the most disturbed cluster in this MACS sample. All morphology parameters are falling on the non-relaxed category which indicates SNR cluster state. In this sample, MACSJ2243.3−0935 is the only cluster which falls in this state. Its position in Fig. 9 is in region 3, i.e. clusters with radio halo region. AM12 classified it as a highly disturbed cluster based on the large separation between X-ray peak or centroid and BCG. Its temperature map [Fig. 7 (right)] is irregular and shows substructures. There are two hot regions, one towards the north-west and another towards the south-east directions. There is trail of hot gas (∼7 keV) connecting these two substructures. These substructures are not visible in the corresponding surface brightness images (Fig. 5). In the lower substructure, the disturbed cool core is visible, while the cool core is completely absent in the upper substructure. As

Figure 8. The sample of MACS clusters is shown in morphology parameter planes along with the P15 sample of clusters. The P15 sample clusters are shown according to their classification – circles are ‘strong relaxed’; diamonds are ‘relaxed’; plus signs are ‘non-relaxed’; and crosses are ‘strong non-relaxed’ clusters. The hexagons are clusters that host radio haloes from Giovannini et al. (2009) and known to be merging clusters. The clusters in the current sample are marked as different colour stars: MACSJ0417.5−1154 (green), MACSJ1131.8−1955 (magenta), MACSJ0308.9+2645 (yellow), MACSJ2243.3−0935 (blue), MACSJ2228.5+2036 (cyan), and MACSJ0358.8−2955 (red).
suggested previously, this cluster is located in the core of supercluster SCL2243.3–0935, surrounded by filaments and it has two concentrations of galaxies corresponding to the two higher temperature regions. We found radio halo at 610 MHz which is situated in between these two substructures. MACSJ2243.3–0935 could be the post-merger cluster.

MACSJ2228.5+2036 and MACSJ0358.8–2955. Based on morphology analysis, both of these clusters are NR clusters; however, we have not detected diffuse radio emission in these clusters. Hsu, Ebeling & Richard (2013) have analysed MACSJ0358.8–2955 cluster in detail using the Chandra, Hubble Space Telescope, and the Keck-I telescope. It shows a high-velocity dispersion (∼1440 km s⁻¹) and a complex merger of at least three substructures. Based on X-ray and optical morphology, they have argued that MACSJ0358.8–2955 is an ongoing merger between two components. However, in the morphology parameter planes, MACSJ0358.8–2955 is situated at the edge of NR clusters. We have put stringent upper limits on the radio powers from these two clusters which is ∼11 times below in the $L_X$–$P_{1.4\text{GHz}}$ plot.

### 6 DISCUSSION AND CONCLUSION

There are six high-$z$, massive ($>7 \times 10^{14} M_\odot$) and non-relaxed MACS clusters in our sample. We analysed their GMRT and EVLA data to detect diffuse radio sources. Out of six, we detect diffuse radio sources in four clusters (67 per cent) while remaining two (33 per cent) are ‘off-state’ clusters. It is well established that not all massive and merging clusters harbour diffuse radio sources (Russell et al. 2011). The typical lifetime of a radio halo is ∼Gyr (Brunetti et al. 2009) which is comparable to cluster merging time-scale. These activities of merger temporarily change the X-ray global properties such as the bolometric luminosity (Ricker & Sarazin 2001). Turbulence generation and propagation processes take time to switch-on the radio emission. This turbulence generation activity is independent of its X-ray morphology. In other words, it is possible that either merging cluster has not generated enough turbulence to form a radio halo or turbulence is already dissipated over a timescale ∼Gyr after the merging process and hence the absence of radio halo. In both of these conditions, host cluster’s X-ray morphology would be disturbed.

There is a steepening in the spectrum of the MACSJ0417.5–1154 radio halo at ∼610 MHz. It is possible that the halo is characterized by a rest-frame cut-off frequency at ∼900 MHz [610 MHz × (1+z)]. This cut-off frequency depends on the magnetic fields and acceleration efficiency in the ICM, and on the energy density of the cosmic microwave background radiation (Cassano 2010). Furthermore, it is believed that ∼Gyr after cluster merger turbulence in the ICM becomes weaker and there are strong synchrotron energy losses which produce a cut-off between ∼100 MHz and GHz frequency range (Cassano 2010; Venturi 2011). However, most of the massive clusters which host haloes have cut-off frequencies ≥1.4 GHz, while in the case of MACSJ0417.5–1154 ($M_{500} = 22.1 \times 10^{14} M_\odot$) the cut-off is at <GHz frequency. One possible scenario is that sufficient amount of time has passed since the merger (∼1 Gyr) and turbulence have decayed in the MACSJ0417.5–1154. In such a scenario, due to less efficient re-acceleration, the cut-off frequency shifts to <GHz frequencies. The position of the MACSJ0417.5–1154 in morphology parameter planes where it is situated at the edge of NR clusters, which indicates less energetic merging events, also supports this scenario. In earlier studies, such ultrasteep spectrum radio halos have been detected in A521 (Brunetti et al. 2008), A1914 (Bacchi et al. 2003), and A2255 clusters (Feretti et al. 1997).

We plotted radio powers at 1.4 GHz ($P_{1.4\text{GHz}}$) and the corresponding X-ray luminosities (0.1–2.4 keV) of known haloes in Fig. 10. Since our sample clusters are at high redshifts, we applied k-correction to their radio powers ($P_{1.4\text{GHz}}$). For applying k-correction, we used a dimming factor of $(1+z)^{(1+\alpha)}$. For calculating this, we used $\alpha = -1.72$ for MACSJ0417.5–1154 and $\alpha = -1.37$ for MACSJ1131.8–1955, while for the remaining sample clusters we used an average spectral index of ∼1.3. Clusters MACSJ1131.8–1955 (A1300) and MACSJ0417.5–1154, which host confirmed haloes, and MACSJ0308.9+2645 and MACSJ2243.3–0935, which host candidate haloes, are situated close to the best-fitting line within intrinsic scatter of the $L_X$–$P_{1.4\text{GHz}}$ correlation. MACSJ0417.5–1154 is showing a cut-off in the spectrum at the rest-frame frequency ∼900 MHz, so obvious that $P_{1.4\text{GHz}}$ is lower than its expected value. From the best-fitting line it is lower by a factor of ∼11. If we use the flux density value from the NVSS observation, then still it is lower by a factor of ∼10. The upper limit on the radio power of the MACSJ0358.8–2955 and MACSJ2228.5+2036 haloes is lower by a factor of ∼11 compared to that expected from the best-fitting line. There could be an alternative to the primary model that secondary CR electrons occur from CR protons’ collisions with ambient thermal protons in a cluster volume. According to Brunetti & Lazarian (2011), the relativistic electrons originating from hadronic (p–p) collisions could produce radio emission which is a factor of ∼10 below the recent
Figure 10. Plot showing the $L_X$ versus $P_{1.4\, \text{GHz}}$ correlation for radio haloes. The data from the current study are labelled with the respective cluster names. The remaining data are from the literature (Kale et al. 2015). The red filled circles are confirmed haloes while the blue squares are candidate haloes. The arrows indicate upper limits. The solid line represents the best fit for the $L_X-P_{1.4\, \text{GHz}}$ relation for haloes: $\log(P_{1.4\, \text{GHz}}) = A \times \log(L_X) + B$, where $A = 2.24 \pm 0.28$ and $B = -76.41 \pm 12.65$. 

$P_{1.4\, \text{GHz}}$ versus $L_X$ trend. Furthermore, recent upper limits on radio luminosity of off-state (non-detection) clusters based on GMRT observations constrain the ratio of CR protons and thermal energy density which is below several per cent and magnetic fields of $\sim$ few $\mu$G (Brunetti et al. 2007; Brown et al. 2011). The latter two clusters in our sample are also reaching to levels at which expected radio emission arise from secondary electrons produced in relativistic $p$–$p$ collisions.

Future low-frequency radio observations are crucial to shed new light on our current understanding of radio halos. In coming years, LOFAR, SKA1-LOW, and MWA will increase the possibilities to utilize the frequency window at the bottom end of the radio spectrum, i.e. below $\sim 200$ MHz. LOFAR will survey the sky at low frequencies with unprecedented depth and sensitivity to detect the faint and ultrafaint diffuse radio sources. Future SKA1 survey will provide us arcsec resolution with excellent surface brightness sensitivity at low frequencies. These future radio observations will permit us to study the formation and evolution of radio haloes in an entirely new range of redshifts and cluster masses. Further improved sensitivities of these newer observations will have implications on the hadronic model.

ACKNOWLEDGEMENTS

We thank the anonymous referee for his/her critical comments that help us to improve this paper. RK is supported through the INSPIRE Faculty Award of the Department of Science and Technology, Government of India. We thank the staff of the GMRT who have made these observations possible. GMRT is run by the National Centre for Radio Astrophysics of the Tata Institute of Fundamental Research. The National Radio Astronomy Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc. The scientific results reported in this article are based in part on data obtained from the Chandra Data Archive and published previously in cited articles.

REFERENCES

Andersson K., Peterson J. R., Madejski G., 2007, ApJ, 670, 1010
Andersson K., Peterson J. R., Madejski G., Goobar A., 2009, ApJ, 696, 1029
Baars J. W. M., Genzel R., Pauliny-Toth I. I. K., Witzel A., 1977, A&A, 61, 99
Bacchi M., Feretti L., Giovannini G., Govoni F., 2003, A&A, 400, 465
Bagchi J., Durret F., Neto G. B. L., Paul S., 2006, Sci, 314, 791
Blasi P., Colafrancesco S., 1999, Astropart. Phys., 12, 169
Böhringer H. et al., 2010, A&A, 514, A32
Bonafede A. et al., 2009, A&A, 503, 707
Bonafede A. et al., 2012, MNRAS, 426, 40
Brown S., Emerick A., Rudnick L., Brunetti G., 2011, ApJ, 740, L28
Brunetti G., Lazarian A., 2011, MNRAS, 410, 127
Brunetti G., Setti G., Feretti L., Giovannini G., 2001, MNRAS, 320, 365
Brunetti G., Venturi T., Dallacasa D., Cassano R., Dolag K., Giacintucci S., Setti G., 2007, ApJ, 670, L5
Brunetti G. et al., 2008, Nat, 455, 944
Brunetti G., Cassano R., Dolag K., Setti G., 2009, A&A, 507, 661
Buote D. A., 2001, ApJ, 553, L15
Buote D. A., Tsai J. C., 1995, ApJ, 452, 522
Buote D. A., Tsai J. C., 1996, ApJ, 458, 27
Cantwell T. M., Scaife A. M. M., Oozeer N., Wen Z. L., Han J. H., 2016, MNRAS, 458, 1803
Carilli C. L., Taylor G. B., 2002, ARA&A, 40, 319
Cassano R., 2010, A&A, 517, A10
Cassano R., Brunetti G., Setti G., 2006, MNRAS, 369, 1577
Cassano R., Brunetti G., Setti G., Govoni F., Dolag K., 2007, MNRAS, 378, 1565
Cassano R., Ettori S., Giacintucci S., Brunetti G., Markevitch M., Venturi T., Gitti M., 2010, ApJ, 721, L82
Cassano R., Brunetti G., Venturi T., 2011, J. Astrophys. Astron., 32, 519
Cassano R. et al., 2013, ApJ, 777, 141
Cavagnolo K. W., Donahue M., Voit G. M., Sun M., 2008, ApJ, 682, 821
Colafrancesco S., Mele B., 2001, ApJ, 562, 24
Condon J. J., Cotton W. D., Greisen E. W., Yin Q. F., Perley R. A., Taylor G. B., Broderick J. J., 1998, AJ, 115, 1693
Conselice C. J., 2003, ApJS, 147, 1
Dickey J. M., Lockman F. J., 1990, ARA&A, 28, 215
Dwarakanath K. S., Malu S., Kale R., 2011, J. Astrophys. Astron., 32, 519
Ebeling H., Edge A. C., Manzi A., Barrett E., Henry J. P., 2001, ApJ, 553, 668
Ebeling H., Edge A. C., Manzi A., Barrett E., Henry J. P., 2001, ApJ, 553, 668
Ensslin T. A., Biermann P. L., Klein U., Kohle S., 1998, A&A, 332, 395
Feretti L., Boehringer H., Giovannini G., Neumann D., 2007, MNRAS, 378, 199
Feretti L., Boehringer H., Giovannini G., Neumann D., 2007, MNRAS, 378, 199
Feretti L., Giovannini G., Govoni F., Murgia M., 2012, A&A Rev., 20, 54
Ferrari C., Govoni F., Schindler S., Bykov A. M., Rephaeli Y., 2008, Space Sci. Rev., 134, 93
Finoguenov A., Böhringer H., Zhang Y.-Y., 2005, A&A, 442, 827
Giacintucci S., 2011, Mem. Soc. Astron. Ital., 82, 541
Govoni F., Bonafede A., Feretti L., Govoni F., Murgia M., Ferrari F., Monti G., 2009, A&A, 507, 1257
Govoni F., Feretti L., 2004, Int. J. Mod. Phys. D, 13, 1549
Govoni F., Feretti L., Giovannini G., Böhringer H., Reiprich T. H., Murgia M., 2001, A&A, 376, 803
Hsu L.-Y., Ebeling H., Richard J., 2013, MNRAS, 429, 833
Intema H. T., 2014, preprint (arXiv:1402.4889)
Intema H. T., van der Tol S., Cotton W. D., Cohen A. S., van Bemmel I. M., Röttgering H. J. A., 2009, A&A, 501, 1185
