A comparison of cosmological filaments catalogues

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ABSTRACT

In this work we compare three catalogues of cosmological filaments identified in the Sloan Digital Sky Survey by means of different algorithms by Tempel et al., Pereyra et al., and Martínez et al. We analyse how different identification techniques determine differences in the filament statistical properties: length, elongation, redshift distribution, and abundance. We find that the statistical properties of the filaments strongly depend on the identification algorithm. We use a volume limited sample of galaxies to characterise other properties of filaments such as: galaxy overdensity, luminosity function of galaxies, mean galaxy luminosity, filament luminosity, and the overdensity profile of galaxies around filaments. In general, we find that these properties primarily depended on filament length. Shorter filaments have larger overdensities, are more populated by red galaxies, and have better defined galaxy overdensity profiles, than longer filaments. Concluding that galaxies belonging to filaments have characteristic signatures depending on the identification algorithm used.

Key words: Cosmology: observations Cosmology: large-scale structure of Universe Galaxies: statistics methods: observational

1 INTRODUCTION

Structures in the Universe are organised hierarchically. From the large scales that are not even bound to small scales as galaxies that act as the cosmological building blocks, all the systems are related. The distribution of matter is a complex web-like structure that includes galaxy clusters, super-clusters, voids, walls and filaments. The filaments are long shaped structures of matter with intermediate densities, galaxy clusters are round shaped with higher densities (above a critical density $\rho_{\text{crit}}$), walls are planar structures, more extended than filaments and with overall lower densities. These structures are entangled, they connect the high density peaks, and leave the majority of the volume empty with the so called voids (Bond et al. 1996) (with typically densities under $0.8\rho_{\text{crit}}$), which tend to become rounder as they evolve. It is not trivial to define precisely, and unequivocally, any of these structures. This is due to their intrinsic complexity, wide range of densities, and diffuse barriers between them. The definition of these structures depends strongly on whether they are to be identified in simulations or in observations. One way to characterise the differences between these structures is to consider clusters, filaments, and walls, as regions with maximal matter density, but with different number of positive eigenvalues of the Hessian matrix of the density field. For the case of clusters, the eigenvalues are all negative, for walls two of them are positive, and for filaments just one (e.g. Chen et al. 2015).

There have been different ways to define filaments in the literature. According to Pogosyan et al. (2009) these objects can be thought as matter ridges that have high densities at the extremes, and a saddle point near the centre. They could also have substructures and then high density peaks near the centre, as discussed by Cantun et al. (2014). Some algorithms to detect filaments take this definition into account. Other authors define filaments as structures that connect massive halos and are traced by other, lesser massive, halos, disregarding the existence of a saddle point between them (e.g. Alpaslan et al. 2014; Park & Lee 2009). There are works that are based on visual inspection (Akamatsu et al. 2017), identification of regions inter-clusters (Pimbblet et al. 2004; Colberg et al. 2005), or on a random configuration of galaxies based on a modelled probability distribution (Tempel et al. 2014a). There are even post-processing algorithms that look for iso-densities surfaces of gas in simulations (Gheller et al. 2015).

Another example in the literature is DisPerSE (Sousbie 2011) a scale-free and parameter-free method to detect
nodes, filaments and walls. DisPerSE identifies structures as components of the Morse-Smale complex of an input density field calculated with the Delaunay tessellation field estimator (DTFE) (Schaap & van de Weygaert 2000). This method proved that it is flexible and robust with both cosmological simulations and observational catalogues. For example Kleiner et al. (2017) use this algorithm to find filaments in the b6F (Jones et al. 2004, 2009) covering the entire southern sky and Kraljic et al. (2018) in GAMA survey (Driver et al. 2009, 2011; Hopkins et al. 2013; Liske et al. 2015). It should be noted that not all algorithms can identify walls explicitly, those methods that do not detect walls may encounter difficulties when dealing with these objects, for instance, by disarming walls as different filaments.

Despite the various drawbacks that the identification of cosmological filaments possess, the study and characterisation of filaments are interesting for different reasons. Several works suggest that in these regions there are unique astrophysical phenomena taking place, corroborating the idea that their nature is related but independent of other structures (Smith et al. 2012).

Tempel & Libeskind (2013) investigate the alignment of spiral/elliptical galaxies, they find that the minor axis of elliptical galaxies tends to be preferentially perpendicular to the hosting filament’s axis, meanwhile bright spiral galaxies tend to be aligned with the host filament’s axis. Zhang et al. (2015) study the alignment between the spin axis of spiral galaxies and the filament direction. They show that the spin axis of spiral galaxies in filaments tends to be preferentially perpendicular to the direction of filaments. The lack of consensus on the observation side regarding galaxy-filament alignments may be due, among other effects, to differences in the definition and method of identification, as well as to unwanted artefacts in sample selection.

Laigle et al. (2015) describe the velocity field around filamentary structures, proposing that this velocity field can be explained with Zeldovich approximation at the saddle points, relating them with the orientation and shapes of galaxies in these environments. Kraljic et al. (2018) study the alignment between the spin axis of spiral galaxies and the filament direction. They show that the spin axis of spiral galaxies in filaments tends to be preferentially perpendicular to the direction of filaments. The lack of consensus on the observation side regarding galaxy-filament alignments may be due, among other effects, to differences in the definition and method of identification, as well as to unwanted artefacts in sample selection.

In this paper we analyse three filament catalogues built using different types of methods and observational samples. In this section we describe each of them to later on study their differences. We call nodes to the filament extreme points, while we refer to the points that form a filamentary path as path nodes.

2 OBSERVATIONAL SAMPLES

In this paper we analyse three filament catalogues built using different types of methods and observational samples. In this section we describe each of them to later on study their differences. We call nodes to the filament extreme points, while we refer to the points that form a filamentary path as path nodes.

2.1 Minimal spanning tree (P19)

The sample of filaments by Pereyra et al. (2019) is constructed by means of a cosmological filament finding algorithm based on a technique borrowed from graph theory: the minimal spanning tree (MST). The authors adopt the definitions of Barrow et al. (1985) and also previously used by others (e.g. Colberg 2007; Alpaslan et al. 2014). A graph is a set of vertex (centres of dark matter halos for numerical simulations, or galaxies for observations typically), edges (they connect vertices) and weights. A MST is the unique set of edges (if all the weights are different) that efficiently connects all the vertices from the initial graph, without closed cycles, and resulting in a minimal sum of the weighs.

The MST describes mainly the distribution of close neighbours and eventually, is capable of making a diffuse MST when the number of vertices is too high, and therefore unable to characterise properly the large scale structure (Stoica et al. 2005). To prevent these drawbacks, the authors limit the MST to a intermediate density region. This region is identified by a Friends-of-Friends (FoF) algorithm designed specifically for flux-limited galaxy surveys, as described in Merchán & Zandivarez (2005). Pereyra et al. (2019) use a transverse large link length of 1.24 h⁻¹Mpc corresponding to an overdensity of δρ/ρ = 1 and a line-of-sight link length of V₀ = 200 km s⁻¹. In addition, each edge in the graph is weighted by the luminosity of the galaxies at its ends.

By using data from SDSS DR12 (Alam et al. 2015), they build the MST, with all the bright galaxies (M_r < −20.5) included in the intermediate density region as path nodes. Then they extract the different branches of the MST. Pereyra et al. (2019) consider as filaments those branches of the tree that have galaxies brighter than m_r < −21.0 at their extremes. As a result, a set of 47249 filaments is achieved. The final catalogue contains the physical length of the filament, elongation, RMS, number of galaxies that
conform the spine of filament and IDs of the galaxies in it, as well as the spectroscopic and photometric properties of the galaxies which were used to build the filament sample. Hereafter we will refer to this sample of filaments as P19.

2.2 Bisous model (T14)
Temple et al. (2014a) use an algorithm called Bisous model which approximates the filamentary structure as multiple cylinders along the filaments that indicate the probability density of galaxies, and the galaxy distribution as a random sample following this density. To determine this structure of cylinders, they start with a random configuration of cylinders that, step by step, is advanced with a stochastic process based on Markov-chains Monte Carlo. The algorithm advances the cylinder locations and orientations to fit with the galaxy distribution. One of the advantages of this approach is that the process relies on the galaxy positions, and no pre-processing, such as computing a smooth density field, is needed. This method is purely based on the geometrical distribution of galaxies, it does not use information about their luminosity or mass.

The filament catalogue is built using the distribution of galaxies in the spectroscopic galaxy sample of SDSS DR8 (Aihara et al. 2011). The catalogue is obtained after concatenating straight lines that have 0.5 Mpc of length. Resulting node’s positions do not necessarily coincide with galaxies or galaxy cluster positions. The catalogue contains the following information about filaments: length, total luminosity of the filament, number of path nodes, co-moving coordinates of each point. There is complementary information for the galaxies associated to these filaments: id of the nearest filament, id of the nearest filament point, and distance from the nearest filament spine, to name only a few. For a detailed description of the filament catalogue see (Tempel et al. 2014a).

Hereafter we will refer to this sample of filaments as T14.

2.3 Filaments linking groups of galaxies (M16)
Martínez et al. (2016) build a sample of filaments using a galaxy group catalogue from Zandivarez & Martínez (2011) as nodes. Filaments consist of pairs of close galaxy groups that are linked by an overdensity of galaxies (see their Fig. 1). Curved or branched structures are not considered because, by definition, these filaments are straight lines joining the nodes. Only groups with virial masses above the average of the group catalogue log(M_{vir}/h^{-1}M_\odot) > 13.5, and in the redshift range 0.05 ≤ z ≤ 0.15 are used. The authors consider that a pair of group is linked by a filament if: (i) they are separated by less than 1000 km s^{-1} in radial velocity; (ii) their projected distance is less than 10h^{-1} Mpc; (iii) the galaxy overdensity in a cuboid-like region between the nodes that is centred on the barycentre of the pair, is greater than a threshold (see details in Martínez et al. 2016).

As result of the identification, at a given filament, we find the physical properties of the nodes such as separation between groups, number of galaxies members, virial radius, virial mass, velocity dispersion. This filament sample has been constructed to study the environmental effects of filaments upon galaxies that are infalling into groups and is not intended to be complete. Hereafter we will refer to this sample of filaments as M16.

2.4 Galaxy sample
For a fair comparison between the three filament catalogues, we use the same parent galaxy catalogue and search for galaxies in this catalogue that lie in the filaments. We use the catalogue of galaxies by Temple et al. (2017), which was downloaded from the SDSS catalogue Archive Server (CAS1, Eisenstein et al. 2011; Alam et al. 2015). These authors added redshifts originated from the Two-degree Field Galaxy Redshift Survey (Colless et al. 2001, 2003), the Two Micron All Sky Survey Redshift Survey (Huchra et al. 2012), and the Third Reference Catalogue of Bright Galaxies (RC3, de Vaucouleurs et al. 1991; Corwin et al. 1994). See Tempel et al. (2014b) for more details. We have adopted the ModellMag magnitudes corrected by extinction and then applied the offset and the k-correction following the empirical k-correction of O’Mill et al. (2011) at z = 0.1. In order to minimise the inclusion of foreground stars (Collister et al. 2007) we used only galaxies with (g − r) < 3 mag.

The volume-complete set of galaxies is determined setting an upper limit of z ≤ 0.137 and a maximum value of absolute magnitude M_g = −20.43, which is computed assuming H_0 = 100 km s^{-1} Mpc^{-1}, Ω_m = 0.3 and Ω_L = 0.7 cosmology. We are aware of possible star-contamination and fiber collisions of the sample, however we find treatment done by Tempel et al. (2017), and O’Mill et al. (2011), good enough for our purposes. This is the catalogue of tracers that we use in this work.

2.5 Random sample of galaxies
We use a random catalogue of galaxies distributed over the same angular distribution of our galaxy catalogue. It consists of ~ 30,000,000 galaxies with (α, δ) coordinates and redshift, 45 times denser than the real galaxy sample. The random sample is a cloning procedure in which every galaxy in our volume limited sample of galaxies is cloned 45 times by assigning to it a random redshift and angular position, bound to mimic the distributions of redshift and angular coverage of the galaxy sample. For this purpose, we constructed an angular coverage mask of the SDSS DR12 using routines from the software HEALPix^2 package (Górski et al. 2005). This procedure does not induce redshift-colour correlations. The relation redshift vs. magnitude is exactly the same for both real and random catalogues.

3 GENERAL PROPERTIES OF FILAMENTS
As thoroughly explained above, the three filament catalogues were obtained through very different processes. The T14 catalogue is a sample of 15420 filaments, while P19 comprises 8350, and M16 3094. In the Figs. 1 and 2 we show the

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1 http://skyserver.sdss3.org/casjobs/
2 http://healpix.sourceforge.net
filaments from the three catalogues over-plotted, in a redshift slice of $z = 0.005 \pm 0.08$, in the plane of the sky Fig. 1, and in a slice of $\delta = 25^\circ \pm 5^\circ$ in Fig. 2 as a pie in the plane of sight. It can be observed that not all the filaments are present in the catalogues and they are not evenly found in redshift. However, there are regions close to bigger structures that seems to cluster the filaments.

The redshift distribution of the three catalogues is shown in Fig. 3, where we can notice that T14 filaments are substantially closer than P19 and M16, as the peaks are located at $z = 0.08, 0.14$ and 0.12 respectively. This shows that T14 is in agreement with the distribution of structures found in Smith et al. (2012). The possible relation between properties like the length of filament as a function of the redshift (as an algorithm bias) has been explored for the catalogues and we do not find appreciable dependence. However, the percentage of galaxies redder than $g - r = 0.7$ slightly increases when decreasing redshift, with the consequence of higher red fractions in general for T14 filaments.

In what follows, we focus on the projected properties of filaments, therefore, we define a sample set of filaments that are perpendicular to the line of sight, and thus in the proper conditions to be stacked. To do so, we use the vector that matches both extremes $v_{0,N-1} = r_{N-1} - r_0$, where $N$ is the number of path nodes, and $r_0$ and $r_{N-1}$ are the position vectors of the extremes. Then we calculate the cosine of the angle between $v_{0,N-1}$ and $r_0$:

$$\cos(\theta) = \frac{v_{0,N-1} \cdot r_0}{||v_{0,N-1}|| |r_0|} \tag{1}$$

If $90^\circ - \Delta \theta_1 < \theta < 90^\circ + \Delta \theta_1$, we consider that the filament belongs to the set 1 (perpendicular to the view axis), and if $\theta < \Delta \theta_2$ or $180^\circ - \Delta \theta_2$ the filament is in the set 2 (parallel to the view axis), where $\Delta \theta_1$ and $\Delta \theta_2$ are tolerance angles that can be tuned to increase the number of selected filaments or to refine the sample.

In an isotropic and homogeneous Universe, all the directions of the filaments are equally likely, however observational effects such as the Fingers of God (Jackson 1972) determine that the number of detected filaments along the visual axis is usually lesser than expected, in particular M16 catalogue excludes this configuration explicitly.

To enhance the signal of the stacking (the proper method is explained in section 4.1), we require that the filaments have the same shape through the following parameters:

- $RMS$: This parameter measures how further away the path nodes are from the positions that make a straight line between the extremes. Since $v_{0,N-1}$ is the vector from the start to the end of the filament, the root mean square of the position of the path nodes from the straight line is

$$RMS = \sqrt{\frac{1}{N} \sum_{i=0}^{N-1} |r_i - r_0 - v_{0,N-1}(r_i - r_0) \cdot v_{0,N-1}|^2} \tag{2}$$

where $N$ is the number of path nodes.

Note that different sized filaments with the same shape will have different values of $RMS$, to solve this, we normalise them by their length to have a scale-independent parameter. High values indicate that the filament is curved or distorted.
and the closer the value is to zero, the more similar it is to a straight line.

- Elongation: This parameter is used by Pereyra et al. (2019) with their filaments, and it is the ratio between the length of the straight line from the extreme and the total length of the path:

\[
\text{Elong} = \frac{|V_{0,N-1}|}{\sum_{i=0}^{N-2} |r_{i+1} - r_i|}.
\]

By construction, this value is always equal or less than 1, the closer the value to 1, the more similar the path is to a straight line.

The analysis described below was done with the set of filaments perpendicular to the line of sight with the additional following criteria:

- T14 filaments: tolerance angle = 20°, elongation = 0.7, and RMS = 0.06.
- P19 filaments: tolerance angle = 30°, elongation = 0.7, and RMS = 0.14.
- M16 filaments: tolerance angle = 30°.

There is a strong correlation between the parameters just shown (i.e. RMS and elongation), we explore the use of both, to clean the samples from irregular shaped filaments.

The most restrictive is elongation and accounting the results of smaller RMS, therefore we only use the elongation criteria to limit our samples.

The distribution of elongations of T14 filaments is quite different than that of P19 filaments. The former has an elongation distribution with a sharp peak very near to 1, and the minimum limit of 0.7 almost does not change the filament set. On the other hand, P19 distribution of elongations is wide, ranging from 0.5 to 1, and the minimum limit of 0.7 considers the ≈ 86% of filaments. We filtered out all filaments with angular size above 4 degrees to avoid contamination from close filaments. As we use a galaxy sample limited to a maximum co-moving distance of 350 Mpc, to avoid doing statistics with filaments above this limit, we filtered them by leaving the sample of filaments closer to this value.

3.1 Filament length

We study global properties of filaments choosing the galaxies inside a region around its axis. This region could be thought as the combination of \(N - 2\) balls of radius \(R\) centred at the inner path nodes, connected by \(N - 1\) cylinders of the same radius and subtracted the balls of the same radius at the end and start of the filament likely related to groups or galaxy clusters. We define the filament radius \(R\) as half its length for the ones shorter than 15 Mpc. For those longer than this limit, we fix \(R\) to 7.5 Mpc. The volume of each filament was estimated with \(V = \pi R^2 l\) where \(l\) is the filament length, although this is not the exact formula for the filament volume, it is a good estimation since we selected straight filaments to study. All galaxies in the region define the properties of the filament, for example the total luminosity. It is worth noticing that according to this membership definition, it is possible that a galaxy could be assigned to more than one filament. The algorithms were applied to both, real and random galaxies, to account how different a filamentary region is to a random distribution of galaxies, according to each catalogue. This helps us to understand what galaxy overdensities the different algorithms are able to find.

As shown in Fig. 4, the bulk of filaments in all catalogues is located between the 5 – 10 Mpc range. However, M16 filaments are limited by 12 Mpc, and P19 filaments extend this limit to 15 Mpc with few exceptions larger than 25 Mpc. These different ranges are explained by the fact that different algorithms (as well as different definitions of filaments) are being used.

3.2 Galaxy overdensity in filaments

In this subsection we analyse the galaxy overdensity in filaments which we compute as:

\[
S_{\text{real}} = \frac{N_{\text{real}}}{N_{\text{rand1}}} \quad (4)
\]

\[
S_{\text{rand}} = \frac{N_{\text{rand2}}}{N_{\text{rand1}}} \quad (5)
\]

where \(N_{\text{real}}\) is the number of real galaxies around filaments, \(N_{\text{rand1}}\) and \(N_{\text{rand2}}\) are the number of galaxies around these objects but instead, from 2 independent random galaxy catalogues.

The overdensity was calculated to both a real and random samples to estimate their distribution functions and how overlapped they were. The regions where the algorithms detect filaments are over-dense, therefore the values for this parameter are above 1 for all the catalogues. But they may differ naturally as a consequence of the different algorithms, for example the P19 algorithm uses a FoF algorithm to discard the low density regions, before actually detecting the filaments, while M16’s defines a filament that is formally the line that joins two groups through an over-dense region. These then are mechanisms that indirectly increase the overdensity of the filaments.

As shown in Fig. 5, the overdensity is approximately 10 for the filaments in the catalogues of P19 and M16, while for
the T14 catalogue it is near 5, compared to their respective random overdensities. This means that these regions contain from 5 to 10 times the amount of galaxies compared to the amount they would have with random galaxies.

Furthermore, the overdensity tends to decrease for long filaments for the catalogues T14 and P19, the correlation is difficult to see in the case of the catalogue M16 due to the limited length range. This could mean that long filaments do actually have less density because short filaments are nearer to galaxy clusters, or, alternatively, that the radius of those filaments is so large that considers regions close to the filament that are not related to the filament itself, thus lowering the overdensity. If this is the case, another radius as a function of the length has to be considered.

3.3 Filament luminosity

The histograms of luminosity of each catalogue can be seen in Fig. 6. It is worth noticing that the three catalogues find filaments in regions with higher luminosity than average. The greatest difference between the random and real galaxies, are seen in M16 and P19 samples, with similar shapes in the distribution and a clear difference in the average luminosity between them. For the T14 sample, both distributions overlap, random galaxies have a broader distribution and real galaxies have a sharper peak. The luminosities are about $10^{11.41} L_\odot$ for real galaxies, and $10^{10.93} L_\odot$ for random galaxies, for T14 filaments. For the P19 sample the quantities are $10^{11.22} L_\odot$ for real galaxies, and $10^{10.31} L_\odot$ for random galaxies. Finally, for the M16 sample, the values are $10^{11.55} L_\odot$ for real galaxies, and $10^{10.56} L_\odot$ for random galaxies. A similar tendency is found in the galaxy number’s histogram (not shown), however there is a higher difference in the distribution for M16 and P19 and quite similar for T14. The average values of the samples are $11.9$ galaxies in T14 filaments, 7.23 for P19 and 15.1 for M16.

3.4 Mean galaxy luminosity in filaments

Another parameter we focus on is the average galaxy luminosity in filaments as the mean of the individual galaxy luminosities in the filament (see Fig. 7). It is not possible to find a correlation between the length and the average luminosity in general, not even a strong difference between the random and real samples. However, the difference is less significant in Temple’s filaments, being a little more luminous the real galaxies (compared with the random samples) for the other two catalogues. For P19 case this is not surprising, because there always exists a path of luminous galaxies between the extremes. In general we see a larger dispersion in the luminosities for short filaments, this is most likely by the effect of low number statistics, as the effect can be reproduced with the random galaxies.

Figure 5. Relation between length and overdensity of filaments for the three catalogues, in white to blue real galaxies and in white to red random galaxies. On the left: filaments from P19 catalogue. Centre: filaments from M16’s catalogue. Right: filaments from the T14 catalogue.

Figure 6. The distributions of galaxy luminosity for real and random galaxies, in colours green and red, respectively. On the left: filaments of the P19 catalogue. In the centre: filaments of the M16 catalogue. Right: filaments of the T14 catalogue.

Figure 7. Relation between length and average galaxy luminosity of filaments for the three catalogues, in white to blue real galaxies and in white to red random galaxies. On the left: filaments from the P19 catalogue. Centre: filaments from M16’s catalogue. Right: filaments from the T14 sample.
3.5 The luminosity function of galaxies in filaments

We compute the luminosity function (hereafter LF) of galaxies in subsamples of filaments defined by their length. Since in this work we use volume limited samples of galaxies our computations of the LFs are restricted to the absolute magnitude range $-23.0 \leq M_r \leq -20.5$, i.e., our LFs are probing only the bright end of the LF. Therefore, when we compute best fit parameters of the Schechter (1976) function below, the $\alpha$ parameter is related to the shape of the bright end of the LF, and does not measure the faint end slope of the LF as is the usual case in the literature. We restrict the LF computations to galaxies in our tracer sample in the redshift range common to the three filament catalogues: $0.05 \leq z \leq 0.137$.

For galaxies in filaments we consider the overall LF of the three samples, and also subsamples of filaments of different redshift space length: $2 - 6$, and $6 - 10$ Mpc. We also compute the LF in groups, and in the field, to compare with the samples of filaments. We construct a sample of galaxies in groups by identifying galaxy groups and clusters using a modified FoF algorithm as described in Merchán & Zandivarez (2005) with a transverse linking length corresponding to an overdensity of $\delta_{\rho}/\rho = 200$ and a line-of-sight linking length of $V_0 = 200$ km s$^{-1}$. We restrict our analysis to massive, $\log(M_{\text{vir}}/M_\odot) \geq 13.5$, groups, i.e., those used in the construction of the M16 filament sample. Our sample of field galaxies comprises all galaxies in the volume under study that are not assigned to groups by the FoF algorithm, nor to filaments by any of the filament samples we use.

We use two standard methods to compute the LF: the step-wise maximum likelihood (SWML, Estathiou et al. 1988) to produce binned LF, and the parametric STY estimator (Sandage et al. 1979) to compute the best-fit Schechter function parameters. We do not attempt to compute the normalisation of the LF, our interest is to study differences in the characteristic magnitude and the shape of the LF between the different subsamples of galaxies. We show in Fig. 8 examples of the LFs we compute, along with their corresponding best fit Schechter functions. In all cases we study, the best fit Schechter model is a good description of the binned LF, therefore, we focus our discussion in the comparison of the parameters $M^*$, and $\alpha$, across samples.

We quote in Table 1 the resulting best fit Schechter parameters of the different LFs we compute, as obtained using the STY estimator. On the one hand, the LF of galaxies in groups has the brightest $M^*$, and the smallest $\alpha$. On the other hand, the LF of field galaxies has the dimmest $M^*$ of all the samples. In a qualitative agreement with the results of Martínez et al. (2016), we find that the characteristic magnitude $M^*$ of the LF of galaxies in filaments takes, in all cases, an intermediate value between those of field and galaxies in groups, but closer to field values. The $\alpha$ parameters of the samples M16 and T14 are also intermediate between the corresponding parameters of field, and group galaxies. However this is not the case of P19, whose $\alpha$ parameter is the largest in all cases.

In the comparison between the LFs of galaxies in filaments in the three catalogues, there are a few points to consider:

(i) The $\alpha$ parameter of P19 LFs is much larger than those of the other two filament samples in all cases, and seems to be independent of filament size.

(ii) Both parameters of the LFs of M16 and T14 are, for all filament sizes, consistent within error-bars.

(iii) The $M^*$ values of the T14 LFs are systematically brighter than those of P19. A similar conclusion can not be drawn when comparing the $M^*$ values of M16 and P19 given the large error-bars in the M16 parameters, due to the smaller sample. However the tendency is for the M16 sample to have brighter values of $M^*$.

(iv) A tendency of shorter filaments to have brighter $M^*$ is seen, this is, however, statistically significant only for the P19 sample.

Despite the differences in the construction of the samples by T14 and M16, their LFs are similar. P19 filaments, on the other hand, have a distinct LF characterised by a much higher value of $\alpha$. This consistency of P19 filaments to have, in all cases, the largest values of $\alpha$, makes them more unlikely to host galaxies in the bright end than the other two samples. Recall that we are probing the bright end of the LF and $\alpha$ is a measure of the convexity of the LF in this magnitude range, and not a measure of the faint end slope, which our samples of galaxies do not probe.

All the characteristics shown in this section suggest that properties of each catalogue strongly depend in the way they were build. Tempel et al. (2014b) algorithm is based on geometrical and stochastic assumptions of the distribution of galaxies, whereas Martínez et al. (2016) algorithm detects overdensities between galaxy groups and Pereyra et al. (2019) method finds filaments as paths of luminous galaxies.

4 SPATIAL GALAXY DISTRIBUTION

In this section we continue our analysis on the differences between the samples of filaments T14, M16, and P19, by focusing on a number of spatial features of the samples. We
use a spatial stacking scheme that we detail below, and compute an adaptation of the two point correlation function that measures the projected clustering of galaxies along the direction perpendicular to the filament axis.

4.1 Spatial Filament Stacking

We stack data of several filaments to enhance the information that different populations of these objects have, increasing the signal/noise ratio. We proceed as follows: Firstly, we define a base set of filaments from each filament catalogue, in which bend and too short filaments are filtered out (see section 3), resulting in a set of straight filaments. Then, for each filament and its surroundings, we define two plane-of-the-sky projected Cartesian coordinates, one alongside the filament direction and centred in one of its extremes (x), and the other perpendicular to it (y). Since every filament has a different length, we re-scale both coordinates in order to have filament length equal to 1. This results in having all filaments with their starting and ending points at the coordinates (0,0) and (1,0) respectively. With this normalisation each part of the filaments, such as the start and end (associated normally with galaxy groups or clusters), the middle filamentary part, the signal beyond the nodes (associated with correlation between connected filaments) and the rest of the field will be stacked at the same places. If no normalisation is done, short/long and far/close filaments will be stacked with different sizes and the galaxies from different parts of each filament will mix together. This procedure is repeated using galaxies from the random catalogue of galaxies.

To process each field near the filament, the angular length of the filament is measured to select an area large enough to cover all the filament field that will be stacked. To avoid summing all the data along the visual axis, that is uncorrelated and adds noise, galaxies with distances further than 10 Mpc from both the start and end of the filament are not considered. The criteria used to consider filaments are the same of the section 3.

Table 1. Luminosity function of galaxies in the field, in groups and in the three samples of filaments: best fit Schechter’s function parameters to the LFs in the absolute magnitude range -23.0 < M_r < -20.5, obtained using the STY estimator.

|        | N_{gal} | α     | M*   | N_{gal} | α     | M*   | N_{gal} | α     | M*   |
|--------|---------|-------|------|---------|-------|------|---------|-------|------|
| Field  | 102,832 | -0.62 | 20.48| -0.06   | -20.13| 0.03 | -        |       |      |
| Groups | 45,701  | -1.04 | 21.19| -0.04   | -20.58| 0.03 | -        |       |      |
| P19    | 6433    | -0.26 | 20.58| -0.01   | 3287  | -0.26| 20.64   | 2314  | -0.26| 20.54|
| M16    | 5259    | -0.60 | 20.61| 654     | -0.70 | 20.8  | 4429    | -0.5  | 20.6  |
| T14    | 47,946  | -0.77 | 20.66| 5740    | -0.9  | 20.8  | 11,052  | -0.8  | 20.7  |

4.2 Mean galaxy overdensity profile of filaments

We build a one-dimensional profile that measures the overdensity of galaxies as a function of the distance to the filaments’ axis. We consider the region defined by 0.15 < x < 0.85, -1 < y < 1, and define the mean overdensity profile as:

\[
I(|y|) = \frac{R(|y|)}{A(|y|)} - 1,
\]

where \(R(|y|)\) is the number of real galaxies, and \(A(|y|)\) the normalised number of random galaxies, at a distance \(|y|\) from the filament axis. The error-bars of these functions are calculated with the Jackknife method, dividing the sample into N/2 subsamples (where N is the number of elements of the sample), therefore calculating each computation excluding two of the filaments and determining the uncertainties. We expect the overdensity profile to reach a maximum near \(y = 0\), while at large values of \(|y|\) the signal should vanish and \(I(|y|) \approx 0\).

In Fig. 10 we show the overdensity profiles for the three samples of filaments, which are consistent with the stackings shown in Fig. 9. As expected, filaments are over-dense regions defined by the large scale structures. In general, the maximum overdensity is in the same magnitude order for all catalogues. The signal is stronger for M16 and P19 because they were constructed considering physical objects as points that define the filament path. On the other hand, the T14 catalogue reaches a much lower value and its profile decreases steeper to the background than the other catalogues.

We study how the overdensity profile is related to filament length by dividing the samples into three different sets per catalogue separating the filaments by length. The bins in filament length we consider are: 4 ± 1, 8 ± 1, and 12 ± 1 Mpc.
Figure 9. Two-dimensional histogram of the filament stacking. Colour indicates the galaxy count. The filament shape show different features for all catalogues, M16 and P19 filaments have clear node radial distributions and T14 have a strong signal and the peaks at nodes are shifted.

In Fig. 11 we show these profiles and it is possible to observe a general decrease of the mean overdensity when increasing the filament’s length, it is also observed in Fig. 5, this is in agreement with the idea of that strong filaments tend to be short bridges that match close galaxy clusters, while for further clusters, filaments are weaker in general (Bond et al. 1996). We find that there is a strong variation for M16. The other catalogues show a milder trend. This fact can be justified again in the way the filaments were constructed. As M16 catalogue is constructed from groups, is understandable that they have a larger correlation for shorter filaments.

Figure 10. Mean overdensity profile of galaxies as a function of the normalised distance to the filaments’ axis, as defined in Eq. 6. Colours represent different samples of filaments as shown in the inset plot.

Figure 11. Mean galaxy overdensity profile around filaments split in three bins of filament length, as indicated on top of each panel. Note that short filaments have a much higher overdensity profile than long ones.

4.3 The overdensity profile of blue and red galaxies

It is possible to use different types of galaxies as tracers of the overdensity profile to study the properties of the filaments. If we separate by colour, it is expected that the red galaxies have higher overdensities at the centre of the filaments, in contrast with the blue galaxies that tend to locate around the filaments as has been shown by Kraljic et al. (2018), and similarly in the works of Dressler (1980) and Blanton et al. (2005) with galaxy clusters. We also expect that the distributions of red and blue galaxies are different whether they are close to galaxy clusters or groups (short filaments) or far away from them (long filaments). Having the
former a higher overdensity of red galaxies at the filament axis compared to that of blue galaxies.

We follow O’Mill et al. (2011) and Fernández Lorenzo et al. (2012), and define red galaxies as those that satisfy $g - r \geq 0.7$, and otherwise for the blue galaxies. This colour separation divides the red-blue bi-modal distribution through the green valley. Blue galaxies comprise the 57.8% of the volume-complete sample, this proportion is roughly constant with $z$ with a slight tendency for lower redshift galaxies to be redder.

The results are shown in Fig. 12 for filaments of 4 Mpc, in the three catalogues. In contrast with Fig. 11, in this figure we use physical units in distance, since the filaments considered here are similar in length. Red galaxies show a high overdensity in the centre of the filaments. Blue galaxies, on the other hand, still have a high overdensity in the centre, but it is lower than the that of red galaxies. This is in good agreement with the works of Dressler (1980) and Blanton et al. (2005) with galaxy clusters. P19 filaments have a higher overall overdensity, M16’s follow slightly below and T14 have the least. This could be explained by the fact that T14 uses a stochastic geometrical algorithm. For filaments of 8 Mpc (Fig. 13), the overdensity decreases, and there is still a clear difference between the overdensity profiles of red and blue galaxies in the samples P19 and M16. T14 filaments do not show a strong difference between red and blue galaxies. Filaments of 12 Mpc have overdensity profiles similar to those of 8 Mpc (Fig. 14). P19 and T14 samples still maintain a slightly higher overdensity for red galaxies. In T14 case this tendency is reverted beyond $\sim 4$ Mpc and M16 filaments do not show a clear signal and they are noisy, most likely because there are few with these lengths.

In general, all catalogues show differences between red and blue galaxies, and, at fixed distance, the overdensity profiles decrease with increasing filament length, regardless of the filament sample. T14 filaments show the least difference between red and blue galaxies, while the largest differences are seen in the M16 case, although for this sample the profiles are noisy when we consider larger filaments. In the P19 case, the smooth overdensity profile is still present for large filaments (12 Mpc), however, the difference between red and blue galaxies vanishes. This can be understood as a consequence of the filament identification method itself, because P19 filaments are constructed through a luminous galaxy path, and therefore confusing red and blue in the outskirts of those galaxies. The increase of uncertainties for the largest set of filaments may be due not only to the small number of filaments, but also to the internal substructure, that will tend to erase the difference in the relative abundance of red and blue galaxies towards the centre of the filaments.

Kraljic et al. (2018) consider filaments as highways of galaxies that can perturb their evolution. If this were the case, galaxies near the nodes should have been flowing...
through the filament for longer time than galaxies in the centre or saddle point. This would cause that the closer a galaxy is to the nodes (as shown by Martínez et al. 2016, and Salerno et al. 2019), the redder it is. We show in Fig. 15 the fraction of red galaxies as a function of the distance to the filaments’ axis. As we move from shorter to larger filaments, the fraction of red galaxies as a function of distance is lower. Shorter filaments are expected to reside in relatively over-dense regions, they are expected to have preferentially red galaxies. Furthermore, it is expected that this short filaments are less independent of the nodes (i.e. behaving like a bridge between them), than larger filaments (Guo et al. 2015). For the largest filaments we analyse, this fraction becomes noisy. Low number statistics do not allow us to study the fraction of red galaxies, as in Fig. 15, but distinguishing between those that are closer to the nodes, or to the saddle points, i.e., binning in the x coordinate.

In Fig. 16, we show the dependence of the red fraction in galaxies in filaments with the filament’s length. We see a general trend that larger filaments have less number of red fraction, tending towards the mean value. That has sense because of the overdensity values of the filaments increases for smaller filaments. In Temple’s case there is a lot of dispersion, and for the stochasticity of their method not bit trend is found.

4.4 Colour vs. Luminosity

Galaxy properties are not independent of each other, and they are affected by environment. As we move from lower to higher density, galaxies are more likely to be brighter, redder, to have earlier morphology, and lower star formation rate. Regarding the broad-band photometrical properties of galaxies, it has been shown by Blanton et al. (2005) that colour is the property that correlates best with local density. In systems of galaxies, colour is also the best predictor of both: the system mass, and the distance to the centre of the system (Martínez & Muriel 2006; Martínez et al. 2008). We have shown in the previous subsection how colour correlates with the distance to the filament’s axis, having red galaxies a higher overdensity profile compared to blue ones. In this subsection we explore how absolute magnitude is related to the distance to filament’s axis, and compare the resulting overdensity profiles with those in the previous subsection, addressing the question of which property is a better predictor of distance.

We focus our study in our samples of filaments with lengths in the range $3 - 5 \, h^{-1}\text{Mpc}$, and consider three absolute magnitude bins, centred in $M_r - 5 \log(h) = -20.75, -21.25, \text{and} -21.75$, and width 0.5. We then compute the overdensity profiles of blue and red galaxies with absolute magnitudes within these ranges. Resulting profiles are shown in Fig. 17. Since our overdensity profiles are a particular type of two-point correlation functions, we fit a power-law of the form $I(|y|) = A|y|^{-\gamma},$ to all profiles in Fig. 17, which we show as dashed lines. Resulting amplitudes and power-law indexes are quoted in Table 2. We find that power-law is a good fit in most cases considered, as can be seen from Fig. 17. Both parameters show, in general, a larger variation from blue to red galaxies at fixed absolute magnitude, than as a function of luminosity at a fixed colour type. For both, blue and red galaxies, and at fixed absolute magnitude, amplitudes are in agreement with the results discussed in the previous subsection. Sorting according to increasing values of the amplitude we have the filaments by T14, P19, and M16. There is no such a clear ranking when we compare power-law indexes, but in general it would be M16, P19 and T14, for increasing $\gamma$, which is in qualitative agreement with Figs. 10 – 14. A singular systematic feature of P19 filaments is that both, the largest amplitude, and the largest power-law index, are obtained for galaxies with $M_r - 5 \log(h) \sim -21.25$ which are the galaxies that P19 use to define their filaments.

Similar results are obtained for longer filaments. We do not include them here in order to not extend further the paper. The main conclusion of this subsection is that the overdensity around filaments is better traced by colour than luminosity. We recall again that in this paper we are probing the bright end of the galaxy population, thus our results concern only to these bright galaxies.

5 CONCLUSIONS

In this paper we present a comparison between different catalogues of cosmological filaments identified by different methods: Pereyra et al. (2019), Tempel et al. (2014b) and Martínez et al. (2016). They differ notably in the way filaments are defined. T14’s algorithm is different than the others in the sense that it uses geometrical assumptions, while the remaining algorithms take in account physical information from the galaxies. This leads to significant differences between those kinds of algorithms that are reflected in the properties of the catalogues.

It is important to note that since these algorithms do not detect walls, some detected filaments could be rather part of walls than a filament itself. This has to be taken into account because walls are different objects and there are other phenomena occurring in them.

The different algorithms do not find the same filaments, instead they find them in common dense regions and in dif-
Figure 16. Fraction of red galaxies in filaments for all the catalogues. It is possible to notice a dependence with the filament length, longer filaments have less fraction of red galaxies.

| Colour | $M_r$     | $A$   | $\gamma$ | $A$   | $\gamma$ | $A$   | $\gamma$ |
|--------|----------|-------|----------|-------|----------|-------|----------|
| Blue   | $-20.75 \pm 0.25$ | $1.90 \pm 0.04$ | $0.84 \pm 0.02$ | $4.1 \pm 0.1$ | $0.54 \pm 0.02$ | $3.35 \pm 0.02$ | $0.71 \pm 0.01$ |
|        | $-21.25 \pm 0.25$ | $2.20 \pm 0.06$ | $0.87 \pm 0.02$ | $4.7 \pm 0.2$ | $0.51 \pm 0.01$ | $3.6 \pm 0.1$ | $0.72 \pm 0.01$ |
|        | $-21.75 \pm 0.25$ | $1.97 \pm 0.05$ | $0.68 \pm 0.02$ | $4.4 \pm 0.4$ | $0.24 \pm 0.03$ | $1.9 \pm 0.3$ | $0.2 \pm 0.1$ |
| Red    | $-20.75 \pm 0.25$ | $2.34 \pm 0.01$ | $0.76 \pm 0.01$ | $10.0 \pm 0.2$ | $0.72 \pm 0.01$ | $4.68 \pm 0.01$ | $0.75 \pm 0.01$ |
|        | $-21.25 \pm 0.25$ | $2.38 \pm 0.07$ | $1.01 \pm 0.03$ | $8.6 \pm 0.2$ | $0.68 \pm 0.01$ | $5.38 \pm 0.02$ | $0.89 \pm 0.01$ |
|        | $-21.75 \pm 0.25$ | $2.50 \pm 0.01$ | $0.96 \pm 0.01$ | $8 \pm 2$ | $0.69 \pm 0.04$ | $2.97 \pm 0.04$ | $0.69 \pm 0.01$ |

Table 2. Best fit power-law parameters to the galaxy overdensity profiles of Fig. 17.

Figure 17. Best fit power-law parameters to the galaxy overdensity profiles of Fig. 17.
end luminosity function. The luminosity function does not vary much with filament length, but there is a tendency of shorter filaments to have brighter characteristic magnitude. Overall, the luminosity functions of T14 and M16 filaments are consistent within errors.

We also develop an statistical tool based on a stacking method that allows us to investigate the spatial distribution of galaxies in and around filaments. With this method we show that the filaments from different catalogues, constructed with various methods, when stacked, they look similar to each other and more extended. The red and blue galaxy distributions also differ between catalogues, the ones that use physical information show more variation between both red and blue profiles while with the T14 catalogue the difference is lesser. We also explore the dependence of the overdensity profiles on luminosity. While there are trends with luminosity, they are in general weaker compared to the dependences on colour. Filament length is an important factor: shorter filaments show a higher overdensity of galaxies. The fraction of red galaxies also vary with the filament length for the M16 and P19 catalogues, such dependence is not found with T14. The red fraction dependence with the length and the position along the filament’s axis has been explored too, but the results that we find are too noisy to reach clear conclusions.

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