The clustering of galaxies as a function of their photometrically estimated atomic gas content

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ABSTRACT
We introduce a new photometric estimator of the HI mass fraction ($M_{\text{HI}}/M_*$) in local galaxies, which is a linear combination of four parameters: stellar mass, stellar surface mass density, NUV−r colour and g−i colour gradient. It is calibrated using samples of nearby galaxies (0.025 < z < 0.05) with HI line detections from the GALEX Arecibo SDSS Survey (GASS) and Arecibo Legacy Fast ALFA (ALFALFA) surveys, and it is demonstrated to provide unbiased $M_{\text{HI}}/M_*$ estimates even for HI-rich galaxies. We apply this estimator to a sample of ∼24 000 galaxies from the Sloan Digital Sky Survey (SDSS)/Data Release 7 (DR7) in the same redshift range. We then bin these galaxies by stellar mass and HI mass fraction and compute projected two-point cross-correlation functions with respect to a reference galaxy sample. Results are compared with predictions from current semi-analytic models of galaxy formation. The agreement is good for galaxies with stellar masses larger than $10^{10} M_\odot$, but not for lower mass systems.

We then extend the analysis by studying the bias in the clustering of HI-poor or HI-rich galaxies with respect to galaxies with normal HI content on scales between 100 kpc and ∼5 Mpc. For the HI-deficient population, the strongest bias effects arise when the HI deficiency is defined in comparison to galaxies of the same stellar mass and size. This is not reproduced by the semi-analytic models, where the quenching of star formation in satellites occurs by ‘starvation’ and does not depend on their internal structure. HI-rich galaxies with masses greater than $10^{10} M_\odot$ are found to be antibiased compared to galaxies with ‘normal’ HI content. Interestingly, no such effect is found for lower mass galaxies.

Key words: galaxies: clusters: general – galaxies: distances and redshifts – cosmology: theory – dark matter – large-scale structure of Universe.

1 INTRODUCTION

Over the past decade, large optical spectroscopic surveys such as the 2dF Galaxy Redshift Survey (2dFGRS; Colless et al. 2001) and the Sloan Digital Sky Survey (SDSS; York et al. 2000) have led to a resurgence in studies of the large-scale structure of the Universe traced by galaxies. There are two main applications of such studies: (a) to constrain cosmological parameters such as the matter density of the Universe $\Omega_m$, Hubble parameter $h$, fluctuation amplitude $\sigma_8$ and neutrino mass in conjunction with constraints from other experiments, such as cosmic microwave background (CMB) or Lyman α forest measurements (e.g. Spergel et al. 2003; Tegmark et al. 2004; Eisenstein et al. 2005), (b) to constrain models for the formation and evolution of the galaxy population.

Traditionally, large-scale structure studies that focus on cosmological applications aim to measure the clustering signal on large scales (tens of Mpc or greater). On large scales, the clustering amplitude depends only on the mass of the dark matter haloes that host the galaxies. All galaxies, regardless of mass or type, trace the underlying dark matter density field in a simple linear fashion, so constraints on cosmological parameters are believed to be robust.

In contrast, studies aimed at constraining galaxy formation focus on the clustering signal on scales less than ∼5 Mpc. On these scales,
the clustering amplitude depends on not only the mass of the dark matter haloes in which galaxies are found, but also the location of galaxies within their host haloes (Benson et al. 2000; Peacock & Smith 2000).

In the current paradigm of galaxy formation within a merging hierarchy of dark matter haloes, galaxies form when gas is able to cool, condense and form stars at the centres of dark matter haloes. At a later stage, the galaxy may be accreted into a larger dark matter halo and become a satellite galaxy in a group or a cluster. Gas is no longer supplied to these galaxies and star formation subsequently shuts down over some time-scale (Kaufmann, White & Guiderdoni 1993; Cole et al. 1994). In more recent models, gas is no longer supplied to central galaxies with central supermassive black holes located in dark matter haloes containing a hot gas atmosphere (Bower et al. 2006; Croton et al. 2006).

One important goal in modern galaxy formation is to understand the physics behind these gas-related ‘accretion’ and ‘quenching’ processes in detail, because the time-scales over which they operate and the way in which their efficiencies scale with halo and/or galaxy mass will determine how the galaxy population as a whole evolves as a function of cosmic epoch. Clustering analysis is a powerful tool in this endeavour. In particular, analysis of the cross-correlation between a specific galaxy subpopulation and a larger ‘reference’ sample allows one to maximize the signal-to-noise ratio of the clustering measurement when the size of the subsample is small. This technique has recently been applied to subsamples of narrow-line galaxies with actively accreting black holes in the SDSS to demonstrate that these galaxies are not triggered by mergers and are found preferentially at the centres of their dark matter haloes (Li et al. 2006b, 2008).

The clustering of galaxies as a function of their neutral gas content should in principle yield very interesting constraints on gas accretion and quenching processes in galaxies (e.g. Popping et al. 2009; Kim et al. 2011). Meyer et al. (2007) determined the two-point autocorrelation function (2PCF) of HI-rich galaxies using 4315 galaxies from the H I Parkes All Sky Survey (HIPASS; Zwaan et al. 2005) and found that HI-selected galaxies exhibit weaker clustering than optically selected galaxies of the same luminosity. Recently Passmoor, Cress & Faltenbacher (2011) measured the 2PCF for an early release of the Arecibo Legacy Fast ALFA (ALFALFA; Giovanelli et al. 2005) sample, finding similar results. The Meyer et al. (2007) study also looked at the dependence of clustering on total H I mass, finding it to be weaker than the dependence on both luminosity and on rotation velocity.

Up to now, there has been no attempt to study how clustering depends on H I mass fraction (i.e. $M_{\text{HI}}/M_*$), a quantity that ought to be much more directly related to accretion and quenching processes that affect the gas content of a galaxy, but not its stellar mass. In addition, a power-law form for the correlation function has been assumed in these previous clustering analyses, which means that information about location of the galaxies within their dark matter haloes (alternatively central or satellite galaxy fraction) cannot be recovered. Finally, because available samples are small, it has not been possible to study clustering as a function of H I mass fraction in conjunction with other galaxy parameters, such as stellar mass or stellar surface mass density. In this work we will demonstrate how an approach that combines H I data for a small, but complete sample of 1000 galaxies and optical data for a much larger sample of galaxies from the SDSS Data Release 7 (DR7; Abazajian et al. 2009) can be used to study the influence of dark matter halo mass and environment on the gaseous properties of galaxies.

The GALEX Arecibo SDSS Survey (GASS; Catinnella et al. 2010) is measuring the atomic gas content of a sample of ∼1000 galaxies with redshifts and stellar masses in the ranges of $0.025 < z < 0.05$ and $10^{10} < M_* < 10^{11.5} M_\odot$. Each galaxy is observed until the HI line is detected or until an upper limit of $\sim 0.015$ in the atomic-to-stellar mass ratio is reached. The GASS galaxies are selected from the SDSS spectroscopic and Galaxy Evolution Explorer (GALEX; Martin et al. 2005), so stellar masses, sizes and structural parameters are available from the MPA/JHU data base (http://www.mpa-garching.mpg.de/SDSS/). The scaling relations of the H I mass fraction of the GASS galaxies ($M_{\text{HI}}/M_*$), as a function of global galaxy parameters such as stellar mass $M_*$, surface mass density $\mu_*$, light concentration index $C$ (defined as $R_{90}/R_{50}$, the ratio of the radii enclosing 90 and 50 per cent of the total $r$-band light) and specific star formation rate $SFR/M_*$, are presented in Catinnella et al. (2010, hereafter C10) and Schiminovich et al. (2010).

Following the work of Zhang et al. (2009), C10 defined a gas-fraction ‘plane’ linking H I mass fraction, stellar surface mass density and near-ultraviolet (NUV)–$r$ colour that exhibited a scatter of 0.315 dex in $\log_{10} M_{\text{HI}}/M_*$, considerably tighter than the relation between H I mass fraction and optical/near-infrared colour studied by Kannappan (2004), which had a scatter of $\sim 0.4$ dex. The improvement in scatter indicates that the H I content of a galaxy scales with its physical size as well as with its star formation rate. In subsequent work, Wang et al. (2011) showed that at fixed NUV–$r$ colour and stellar surface density, galaxies with larger H I gas fractions have bluer outer discs.

In this paper, we include the colour gradient of galaxies as an additional parameter in our fits. This produces a relation with similar scatter, but that better predicts the H I mass fraction of the most gas-rich galaxies in our samples. We use this relation to predict the H I content of the galaxies in our SDSS/DR7 sample. We then study how clustering depends on both ‘pseudo’-H I mass fraction and a ‘pseudo’-H I excess/deficiency parameter, which we define as the deviation in the predicted H I content of a galaxy from the average H I content of all galaxies of the same stellar mass and surface mass density. This ‘pseudo’-H I excess/deficiency parameter depends on a combination of NUV–$r$ colour and $g$–$r$ colour gradient. Finally, we compare our results with clustering predictions from the semi-analytic galaxy formation models of Fu et al. (2010, hereafter F10) and Guo et al. (2011, hereafter G11). The main way in which the F10 model differs from the G11 model is that it includes simple prescriptions for molecular gas formation processes.

The motivation behind expressing the results in this paper in terms of ‘pseudo’-H I fraction, rather than in terms of directly measured photometric quantities, is because this provides insight into physical processes regulating the gas supply in galaxies. The semi-analytic models make a host of assumptions about how gas is accreted from the surround dark matter halo and then consumed into stars. By comparing clustering as a function of gas fraction in the models with the data, we hope to ascertain whether these assumptions are correct or whether there are discrepancies that warrant further investigation.

Throughout this paper we have assumed a cosmology with $\Omega = 0.3$, $\Lambda = 0.7$ and $H_0 = 70\,\text{km}\,\text{s}^{-1}\,\text{Mpc}^{-1}$ when computing observed quantities. A Hubble constant of $H_0 = 100\,\text{km}\,\text{s}^{-1}\,\text{Mpc}^{-1}$ is assumed when presenting correlation function measurements. We note that the F10 and G11 models are based on simulations with $\Omega = 0.25$ and $\Lambda = 0.75$. This will make a small difference in the comparison between data and models. We note that the focus of this paper is not on obtaining precision fits to the data, but on identifying major discrepancies that may lead us to change the input physics in the model.

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2 DATA

2.1 GASS galaxy sample

The parent sample of GASS consists of 12 006 galaxies selected from the region of sky where the sixth Data Release (DR6; Adelman-McCarthy et al. 2008) of the SDSS overlaps the maximal ALFALFA footprint. All galaxies are selected to have stellar masses $M_*>10^{10} M_\odot$ and redshifts in the range of $0.025<z<0.05$. The GASS sample is constructed by randomly selecting a subset of ~1000 galaxies from the parent sample within the footprint of the GALEX Medium Imaging Survey so that the distribution in stellar mass is flat. The targets are observed with the Arecibo radio telescope until detected or until an HI mass fraction $M_{HI}/M_*$ limit of 1.5–5 per cent is reached. In this work, we use the representative sample of 480 GASS galaxies, including 293 detections and 187 non-detections described in Catinella et al. (2012). Details of the GASS survey design, target selection and observing procedures can be found in C10.

2.2 SDSS galaxy samples

We have constructed two galaxy samples from the SDSS/DR7.

The first sample, which will serve as our reference sample in the clustering analysis, is a magnitude-limited sample of 66 461 galaxies with $r<17.6$, $-24<M_{i,1}<-16$ and redshifts in the range of $0.025<z<0.05$. Here, $r$ is the r-band Petrosian apparent magnitude, corrected for Galactic extinction, and $M_{i,1}$ is the r-band Petrosian absolute magnitude, corrected for evolution and K-corrected to its value at $z=0.1$. These selection criteria, with the exception of the redshift range, are the same as in our previous papers where we studied the clustering of galaxy luminosity and stellar mass (e.g. Li et al. 2012). We have generated random samples that have the same sky coverage as well as the same position- and redshift-dependent selection effects as the reference sample. Details of this procedure are presented in our previous papers (e.g. Li et al. 2006a).

The second sample contains 36 136 galaxies, and is a subset of the reference galaxies with stellar masses in the range of $10^{9.5} M_\odot < M_* < 10^{11} M_\odot$. In the next section we will estimate an HI mass fraction for each galaxy in this sample using our newly calibrated photometric estimator. We select a number of subsamples binned according to $M_*$ and $M_{HI}/M_*$, and cross-correlate these with both the reference and random samples.

2.3 Physical properties of galaxies

The physical quantities necessary for this work include stellar mass $M_*$, stellar surface mass density $\mu_*$, NUV $-r$ colour and colour gradient $\Delta g_{r-i}$. Stellar masses are derived from SDSS photometry using the methodology described in Salim et al. (2007). These masses are publicly available at http://www.mpa-garching.mpg.de/SDSS/DR7. The stellar surface mass density is defined as $\mu_* = M_*/(2\pi R_{500}^2)$, where $R_{500}$ is the physical radius in units of kpc that contains half the total light in the $z$ band. The NUV magnitude is provided by the GALEX pipeline and the NUV $-r$ colour is corrected for Galactic extinction following Wyder et al. (2007) with $A_{NUV-V_0} = 1.9807 A_r$, where $A_r$ is the extinction in r band derived from the dust maps of Schlegel, Finkbeiner & Davis (1998). The $g-i$ colour gradient is defined as $\Delta c_{g-i} = (g-i)_{\text{int}} - (g-i)_{\text{ext}}$, where $(g-i)_{\text{int}}$ and $(g-i)_{\text{ext}}$ are the $g-i$ colours in the inner and outer regions of the galaxy. The inner region is enclosed by $R_{500}$, the radius containing half the $r$-band light. The outer region is defined as the area between $R_0$ and $R_0$, the radius enclosing 90 per cent of the $r$-band light. A negative value of $\Delta c_{g-i}$ implies that the outer region of the galaxy is bluer than the inner region.

2.4 Semi-analytic model galaxy catalogues and mock SDSS samples

In this paper we compare our observational results to predictions from the galaxy formation models of G11 and F10. Both models were created by implementing simple prescriptions for baryonic astrophysics on merger trees that follow the evolution of the halo/subhalo population in the Millennium Simulation (MS; Springel et al. 2005), a cubic region 500 $h^{-1}$ Mpc on a side with mass resolution $\sim 10^{10} M_\odot$. The G11 model is the most recent semi-analytic model from the Munich group, in which the treatments of supernova feedback, galaxy size, photoionization suppression and environmental effects on satellite galaxies have been significantly updated. G11 demonstrated that their model provided excellent fits not only to the luminosity and stellar mass functions of galaxies derived from SDSS data, but also to recent determinations of the abundance of faint satellite galaxies around the Milky Way. The clustering properties of galaxies as a function of stellar mass predicted by the model are in good agreement with SDSS data for masses above $6 \times 10^{10} M_\odot$ and at separations larger than 2 $Mpc$. On smaller scales, lower mass galaxies are predicted to be substantially more clustered than observed.

The F10 model is based on an earlier version of the Munich semi-analytic code, which is described in detail in Croton et al. (2006) and updated in De Lucia & Blaizot (2007, hereafter DB07). The main new aspect of this model is that galactic discs are represented by a series of concentric rings in order to track the evolution in the gas and stellar surface density profiles of galaxies over cosmic time. In addition, two simple prescriptions for molecular gas formation processes are included: one is based on the analytic calculations by Krumholz, McKee & Tumlinson (2009) and one is a prescription where the $H_2$ fraction is determined by the pressure of the interstellar medium (ISM; Blitz & Rosolowsky 2006). The model is currently being configured to operate on the latest code of G11. The comparison in this paper will be restricted to the model published in the F10 paper and to the $H_2$ formation prescription of Krumholz et al. (2009).

In this paper, the different treatments of gas stripping in the G11 and F10 models are of interest to us. In most semi-analytic models including DB07, hot gas in a halo is assumed to be stripped immediately after the halo has been accreted on to a larger halo. In the G11 model this prescription has been modified. Satellite galaxies that still are attached to a subhalo within the larger virialized ‘parent’ halo are still able to accrete gas. This new treatment was motivated by observational findings and hydrodynamical simulations which revealed that the hot atmosphere of massive satellite galaxies may survive for a considerable time after accretion (see G11 and references therein for details). This change primarily affects satellite galaxies located in the outer regions of their host dark matter haloes. The time-scale for gas to be depleted and star formation to stop becomes significantly longer.

We have constructed a set of 50 mock SDSS galaxy catalogues from the G11 model using both the sky mask and the magnitude and redshift limits of our SDSS reference sample. Detailed description of our methodology can be found in Li et al. (2006b) and Li et al. (2007). These mock catalogues allow us to derive realistic error
estimates for the statistics measured below, including both sampling and cosmic variance uncertainties.

3 ESTIMATING H I MASS FRACTIONS FOR THE SDSS GALAXIES

There have been a number of attempts to calibrate colours (e.g. Kannappan 2004) or emission-line equivalent widths (e.g. Tremonti et al. 2004; Erb et al. 2006; Bouché et al. 2007) as proxies for the gas-to-stellar mass ratio in galaxies. Zhang et al. (2009) proposed a method motivated by the Kennicutt–Schmidt star formation law (Schmidt 1963; Kennicutt 1998) that combines colour and surface brightness to estimate the H I-to-stellar mass ratio. They used a sample of 800 galaxies with H I mass measurements from the HyperLeda catalogue (Paturel et al. 2003) and optical photometry from the SDSS to calibrate a relation linking these quantities. In subsequent work, C10 used an unbiased sample of galaxies with H I measurements from GASS to show that \( \frac{M_{\text{HI}}}{M_*} \) can be well approximated by a linear combination of NUV-to-optical colour (\( NUV - r \)) and stellar surface mass density (\( \mu_* \)) with a 1\( \sigma \) scatter of \( \sim 0.3 \) dex. However, as could be seen from fig. 12 of C10, galaxies detected by the much shallower ALFALFA survey in redshift range that the GASS sample had significantly higher H I mass fractions, and were also systematically displaced from the C10 plane. This result would seem to imply that the H I masses of the most gas-rich galaxies in the local Universe cannot be reliably inferred from their UV/optical properties.

However, a recent study by Wang et al. (2011) focusing on the H I-rich galaxies from the GASS and ALFALFA samples has revealed that unusually H I-rich galaxies have bluer than average outer discs. Motivated by this finding, we now propose an updated photometric estimator that includes both stellar mass and the gradient in colour as additional parameters.

In Fig. 1, we plot some of the galaxy samples we will be working with in this paper in the two-dimensional plane of \( NUV - r \) colour and stellar surface mass density. There are three grid centres (plotted in red squares) located well within the ‘red sequence’ with good mean H I mass measurements and these provide a check on whether our H I mass fraction estimators work well in regime where galaxies are gas-poor on average. As can be seen, the combination of the GASS and ALFALFA data, as well as the stacked results, covers the region of \( NUV - r \) versus \( \log_{10} \mu_* \) parameter space reasonably well. The GASS galaxies and stacked results are offset to somewhat higher values of stellar surface mass density, because these samples are restricted to galaxies with stellar masses larger than \( 10^{10} \) \( M_\odot \).

Our new estimator is defined by

\[
\log_{10} \frac{M_{\text{HI}}}{M_*} = a \log_{10} \mu_* + b(NUV - r) + c \log_{10} \frac{M_*}{M_\odot} + d \Delta_{g-i} + e, \tag{1}
\]

where \( \Delta_{g-i} \) is the colour gradient defined in Section 2.3. The coefficients are determined by minimizing the residuals from the plane using the 293 H I detections in the GASS sample. When carrying out the fit, we weight each galaxy by the mass-dependent selection function of the GASS survey (note that, although C10 did not apply this weighting, the results are entirely consistent). The 1\( \sigma \) scatter in our new estimator is 0.31 dex, very similar to that of the old one. Fig. 2 illustrates how this new estimator improves the H I mass fraction estimates.

In the left-hand panel of the figure, grey dots show the H I plane of C10 for GASS galaxies, while coloured stars show the stacked results. In the right-hand panel of the figure, grey dots show the same C10 plane for a sample of ~7000 galaxies from the a.40 catalogue of the ALFALFA survey (Haynes et al. 2011) with stellar masses above \( 10^8 \) \( M_\odot \) and redshifts below 0.06. As can be seen, the majority of the grey points in the right-hand panel lie above the relation.

The H I plane given by the new estimator in equation (1) is plotted in red open circles or red dots in both panels of Fig. 2. There is rather little change for the majority of galaxies in the GASS sample. However, the H I-rich galaxies in the ALFALFA sample that were previously displaced to higher than predicted H I mass fractions are now mostly located well within the 1\( \sigma \) region of the new relation.

We note that this reduction in the systematic offset for H I-rich galaxies could not be achieved by introducing a single new parameter into the fit (i.e. only \( \Delta_{g-i} \)). Equation (1) implies that the predicted gas fraction scales more strongly with colour gradient in high-mass galaxies than in low-mass galaxies. The most likely reason for this is that massive galaxies have larger bulge-to-disc ratios than less massive galaxies. Fabello et al. (2011) showed that the H I content of a galaxy did not depend on its bulge-to-disc ratio; the H I mass fraction only depended on the properties of the disc. It is thus likely that the H I fraction correlates with the colour gradient of the disc and the bulge is a contaminant when determining the colour gradient. At present, we do not have bulge/disc decompositions for all the galaxies in our sample, so we do not investigate this.

1 Note that the bins on the red sequence have been coloured in green and lie very close to an extrapolation of the best-fitting line through the other bins, indicating that the C10 plane still yields an accurate prediction of mean H I mass fraction for galaxies on the red sequence.
hypothesis in more depth. Another effect that may be important is that massive galaxies contain more dust, and this may change the relation between colour gradient and $H_1$ fraction.

We now test whether our new estimator exhibits any remaining systematic biases by checking whether the residuals are correlated with any intrinsic galaxy property. In Fig. 3 we plot the residuals for the C10 estimator (grey dots) and for the new estimator (colourful dots) as a function of $M_\ast$, $\mu_\ast$, $NUV - r$ and $R_{90}/R_{50}$. We only show results for the ALFALFA sample, where the new estimator does change the $H_1$ mass fraction predictions by a significant amount.

Fig. 3 shows that the new estimator leads to a significant reduction in the large positive residuals for $H_1$ rich galaxies with low masses and stellar surface mass densities, blue colours and low concentration indices. The new estimator does not reduce the residuals for galaxies with high stellar surface densities, red colours and high concentration indices. There is still a subpopulation of such galaxies that are $H_1$-rich and where equation (1) fails to predict the $H_1$ content accurately. An example of such a system is discussed briefly in C10. In addition, one might worry that the $H_1$ mass fraction estimation may be biased in this regime, because many red galaxies are not detected in both the GASS and ALFALFA surveys.

In the left-hand panel of Fig. 2 we plot the results of the $H_1$ stacking analysis by Fabello et al. (2011). By stacking samples of a few hundred galaxies, Fabello et al. (2011) were able to estimate mean $H_1$ mass fractions for galaxies with $NUV - r$ colours in the range of 4–6 (shown as green stars on the plot). As can be seen, the C10 estimator accurately reproduces the stacked results with a $1\sigma$ scatter less than 0.07 dex, even for the reddest stacks. Unfortunately, a similar test is not possible for our new estimator, because the sample of SDSS galaxies with available ALFALFA coverage is too small to carry out a stacking analysis using four different galaxy parameters instead of two. Therefore, in what follows, we divide our galaxies into ‘red’ or ‘blue’ using a mass-dependent colour divider:

$$(NUV - r)_c = 0.5 \log_{10}(M_\ast/M_\odot) - 1.$$  (2)

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Projected two-point cross-correlation function, $w_p(r_p)$, for six subsamples with equal number of galaxies from the GASS, selected by the observed $\text{H} \text{i}$ mass fraction (triangles connected by blue dotted lines) or the predicted value (squares connected by red solid lines). The average value of the $\text{H} \text{i}$ mass fraction of each subsample is indicated. The result for the whole sample is plotted in the black solid line and repeated in every panel for reference.

For ‘red’ galaxies with $\text{NUV} - r > (\text{NUV} - r)_\text{cut}$ we use the old estimator, while the new estimator is applied to ‘blue’ galaxies with $\text{NUV} - r < (\text{NUV} - r)_\text{cut}$.

We now carry out a test to see whether the estimator introduces any systematic bias in galaxy clustering analyses. We divide all the GASS galaxies including non-detections into six subsamples of equal size, using both the measured value of $M_{\text{H} \text{i}}/M_*$ and the predicted value, and we compute the projected two-point cross-correlation functions (2PCCF), $w_p(r_p)$ of these subsamples with the SDSS reference sample. We find that it is very important to take into account the effect of the errors in the predicted value of $M_{\text{H} \text{i}}/M_*$ when comparing the results using the photometric estimator with the results using the real HI measurements. The effect of errors is to weaken the clustering trends as a function of $M_{\text{H} \text{i}}/M_*$, particularly in the tails of the distribution. Here, we model the effect of the errors by adding a random component to the measured value of $M_{\text{H} \text{i}}/M_*$ that follows a Gaussian distribution function with a width of 0.31 dex. Results using the measured HI mass fractions convolved with the Gaussian distribution of errors are shown in blue in Fig. 4, while results using the photometric estimator are shown in red. The errors in the $w_p(r_p)$ measurements are computed using the bootstrap resampling technique. The 2PCCF for the whole GASS sample is plotted as a black solid line in each panel for reference. Fig. 4 shows that the two $w_p(r_p)$ calculations agree well with each other. We have repeated the same analysis, using the measured $\text{H} \text{i}$ mass fraction for the GASS galaxies without including the effect of the errors, and found that the results change very little, indicating that the smearing by the $\text{H} \text{i}$ predictor on the correlation function is sufficiently small.

Although the GASS sample is small, we can still see that both the amplitude and shape of the 2PCCF show strong systematic trends with increasing $\text{H} \text{i}$ mass fraction. $\text{H} \text{i}$-rich galaxies are less strongly clustered on all scales, with more pronounced one-halo to two-halo transitions at $\sim 1$ Mpc. Since galaxy clustering depends on a variety of physical properties, in particular on stellar mass, it is unclear to what extent the effect seen from Fig. 4 is due to $\text{H} \text{i}$ content only. We will address this point in the next section.

4 CLUSTERING AS A FUNCTION OF $M_{\text{H} \text{i}}/M_*$ AND COMPARISONS WITH SEMI-ANALYTIC MODELS

In this section we apply our new photometric estimator to our full SDSS DR7 galaxy sample to study the dependence of clustering on $\text{H} \text{i}$ mass fraction. We compare our results with predictions from the G11 and F10 models. It is well known that clustering depends strongly on galaxy stellar mass, so the analyses are always carried out in narrow mass intervals. In order to take errors in the photometric estimator into account, we convolve the $\text{H} \text{i}$ mass fractions predicted by the models with a Gaussian distribution function of width 0.3 dex in $\log_{10}(M_{\text{H} \text{i}}/M_*)$.

Before we begin, we demonstrate that the models reproduce average trends in $\text{H} \text{i}$ mass fraction as a function of stellar mass reasonably well. In Fig. 5, the black solid curve shows the median value of $\log_{10}M_{\text{H} \text{i}}/M_*$ as a function of stellar mass for galaxies in the GASS survey, while the grey shaded region indicates the 16th to 84th percentile ranges of $\log_{10}M_{\text{H} \text{i}}/M_*$. Note that the galaxies without $\text{H} \text{i}$ line detections are assigned an $\text{H} \text{i}$ mass equal to the upper limit. This is why the black curve and the shaded region do not fall below $\log_{10}M_{\text{H} \text{i}}/M_* \sim -1.82$ (see C10 for details on the detection limits of the survey). We now perform the same analysis for the simulated galaxies and the results are shown in red and blue for the F10 and G11 models, respectively. We see that the G11
model yields a higher median value of H I gas mass fraction at a given value of $M_*$, when compared to both the data and the F10 model.

There are two reasons for this: (1) the G11 model does not account for the partition of the neutral gas into different components. The F10 model includes simple prescriptions for the formation of molecular gas and also properly takes into account the contribution of helium when making predictions for H I content. (2) The F10 model parameters are explicitly adjusted so as to match the H I mass functions determined by existing H I surveys such as HIPASS and ALFALFA. Kauffmann et al. (2012) have shown that the F10 model can also reproduce the distribution of $M_{HI}/M_*$ for the population of galaxies with detectable gas, but the model does not provide a fully accurate description of the population of galaxies without detectable cold gas. We note that Fig. 5 includes both populations, so the fit to the data is not as good as that shown in fig. 2 of the Kauffmann et al. paper. Since the GASS is a survey only for high-mass galaxies with detectable gas, we can of course check whether the models predict a similar $wp(r_p)$ on stellar mass. The agreement with observations is equally good for both models on large scales. The F10 model appears to provide a somewhat better fit to the clustering amplitude on scales below $\sim 1$ Mpc.

In spite of the good agreement as shown above, there still exist discrepancies in some cases. In order to carry out meaningful comparisons between data and models, we order all the galaxies in a given stellar mass range by increasing $M_{HI}/M_*$ and divide the galaxies into 10 subsamples, each containing 10 per cent of the whole sample. We analyse the dependence of $wp(r_p)$ on H I mass fraction as a function of H I mass fraction percentile instead of the absolute value of H I mass fractions. In order to provide a more intuitive feel for our results, we present our measurements in terms of bias factor, defined as the ratio of the $wp(r_p)$ for a given H I-selected subsample to the $wp(r_p)$ of all galaxies in the corresponding stellar mass range. In Fig. 7, we plot this bias factor as a function of percentile in log$_{10} M_{HI}/M_*$, with H I mass fraction increasing from left to right. Results for different intervals in stellar mass are shown in different rows, while results evaluated on different projected scales $r_p$ are shown in different columns. The data are shown in black curves with shaded regions indicating the 1σ errors that are estimated from the bias factor measurements of the 50 mock SDSS catalogues, while the F10 and G11 models are shown in red circles and blue triangles, respectively.

As we will now show, this decrease in bias can be understood in a simple way in terms of an increasing ratio of central-to-satellite galaxies as a function of increasing H I mass fraction. To prove that this is the case, we classify each galaxy in our sample as either a central galaxy or a satellite galaxy based on whether it is more massive than all companions within a cylinder with projected radius $R_{max}$ and a line-of-sight depth of $\pm 1000$ km s$^{-1}$. Here, $R_{max}$ is set to twice the virial radius of the host dark matter halo of the galaxy. We have adopted the stellar mass–halo mass relation derived by Guo et al. (2010) to estimate a halo mass for the galaxy, and then estimate a ‘virial’ radius of the halo using the model of Eke, Navarro & Steinmetz (2001). In addition, we require that a central galaxy should not fall within $R_{max}$ of any other more massive galaxy.

In the right-hand panels in Fig. 7, we plot the fraction of central galaxies, $f_{cen}$, as a function of H I fraction percentile (the black solid line). We see that $f_{cen}$ increases with increasing H I content, with the effect stronger at low stellar masses.

We can of course check whether the models predict a similar increase in $f_{cen}$ as a function of H I gas fraction. In the right-hand column of Fig. 7, the dashed curves show the true values of $f_{cen}$ as a function of H I fraction percentile for the F10 (red) and G11 (blue) models. The central fractions in the F10 model are lower than those in the G11 model at low H I mass fractions, particularly for galaxies with low stellar masses. This reflects the fact that gas consumption times in satellite galaxies are longer in the G11 model than in the F10 model.

In order to make a fair comparison with observations, we have also computed $f_{cen}$ for the model galaxies in exactly the same way as in the observations. In brief, we project the model galaxies on to the $x = y$ plane and take the $z$-axis as the line-of-sight direction (i.e. we adopt the distant observer approximation). We then apply exactly the same procedure described above to classify each galaxy as central.

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2 We note that this only makes sense for the G11 model if the intrinsic scatter in H I-to-H$_2$ ratio in real galaxies does not change the ranking of $M_{HI}/M_*$ with respect to $[M_{HI} + M_{H_2}]/M_*$. Saintonge et al. (2011) show that the average value of $M_{HI}/M_*$ is around 1/3 and the molecular gas mass very rarely exceeds the atomic gas mass, so this is likely to be close to correct.
Figure 6. The projected two-point cross-correlation function \( w_p(r_p) \) for galaxies in bins of stellar mass with respect to a reference sample of all galaxies. Results for the SDSS DR7 sample are shown in black triangles, while those for the Fu et al. (2010) and Guo et al. (2011) models are shown in red solid lines and in blue dashed lines, respectively. Errors on the SDSS results are estimated from a set of 50 mock galaxy surveys that have the same selection effects as the real SDSS sample.

Figure 7. Panels on the left-hand side show the clustering bias factor as a function of percentile in \( \log_{10}(M_{\text{HI}}/M_\ast) \), at different separations (panels from left to right) and for different stellar mass ranges (panels from top to bottom). The solid line shows observed bias factors, while the shaded region indicates the errors on the observed bias factors estimated from the mock catalogues. Model results are shown in red/blue for the F10/G11 models. In the rightmost column, the central galaxy fraction is plotted as a function of percentile in \( \log_{10}(M_{\text{HI}}/M_\ast) \). The black curves with error bars show the central fractions estimated from the data (see text). The dashed red/blue curves show the true central fractions from the models. The solid red/blue curves show the central fractions in the models when estimated in the same way as in the data.

or satellite. The \( z \)-axis peculiar velocities of the galaxies are added to their \( z \)-axis positions. In addition, halo masses and virial radii are not taken from the model catalogue, but are estimated exactly the same way as for the observational data. Results are plotted as solid red and blue curves in the right-hand panels of Fig. 7.

We note that the true central fractions are always smaller than the ones that use a classification technique based on whether or not brighter companions are found in cylinders around the galaxy. However, our classification technique preserves the shape of the relation between central fraction and \( \text{HI} \) gas fraction percentile, as well as the differences between the F10 and G11 models. We also see that the central fractions estimated in cylinders in the simulation agree reasonably well with the data. As was the case for the bias factor, the behaviour of the central fraction as a function of \( \text{HI} \) mass fraction percentile in the models is somewhat different to what is seen in the observations, particularly at low stellar masses.

In summary, the general agreement with the models supports our conjecture that the trends in bias factor as a function of \( \text{HI} \) mass fraction mainly arise as a result of trends in the satellite-to-central ratio.
5 SCALE DEPENDENCE OF THE BIAS FOR GALAXIES WITH EXCESS/DEFICIENT HI CONTENT

In the previous section, we studied how the bias factor changes as a function of normalized HI mass fraction for galaxies in different stellar mass bins. The results presented in Fig. 7 clearly show that the change in the bias factor between ‘gas-poor’ and ‘gas-rich’ galaxies depends on scale $r_p$.

In this section, we analyse the scale dependence of the change in bias factor for both gas-deficient and gas-rich galaxies. We note that Haynes & Giovanelli (1984) defined gas deficiency to be the difference in HI content between cluster galaxies and ‘field’ galaxies of the same morphological type and size. There are some difficulties with this definition, including the definition of ‘field’ and the fact that Hubble classification is problematic in rich clusters. Some more recent analyses have used a type-independent deficiency parameter that compares all galaxies to a fixed HI surface density (Chung et al. 2009). One worry with this is that the mean HI content of galaxies scales strongly with galaxy parameters such as stellar mass and surface density (C10).

In this analysis, we will adopt a flexible approach to defining ‘pseudo’-HI deficiency parameters. We will analyse the photometrically predicted HI content with respect to galaxies of the same stellar mass, with respect to galaxies of the same mass and size, and with respect to galaxies of the same mass, size and colour (as in Cortese et al. 2011).

In the upper panels of Fig. 8, we show the change in bias factor between the 10th and 50th percentile bins in $\log_{10}(M_{\text{HI}}/M_\ast)$. This serves as a test of gas-stripping mechanisms in gas-deficient galaxies. The lower panels show the change in bias factor between the 100th and 50th percentile bins in $\log_{10}(M_{\text{HI}}/M_\ast)$. This serves as a test of gas accretion mechanisms in gas-rich galaxies.

Results for the SDSS DR7 galaxies are shown as black curves in Fig. 8. Grey shaded regions indicate the 1σ errors on our estimates, obtained from the 50 mock SDSS catalogues. Results for the F10 and G11 models are shown as red and blue curves and we plot our results in three different stellar mass ranges. As can be seen from the plot, the change in bias factor between gas-deficient galaxies and galaxies with typical gas fractions is most pronounced for low stellar mass systems. The bias factor difference peaks at relatively small physical scales ($\sim$100–300 kpc). For the most massive galaxies with $9.5 < \log_{10}(M_\ast/M_\odot) < 11$, there is little change in bias on any scale. For galaxies with $9.5 < \log_{10}(M_\ast/M_\odot) < 10$, the increase in clustering amplitude from galaxies with typical gas mass fractions to the most gas-deficient objects reaches a factor of 2 on scales of a few hundred kpc. On scales larger than 2–3 Mpc, there is no significant change in clustering amplitude. The results are consistent with the idea that gas quenching is driven by processes that are internal to dark matter haloes. The models agree well with the data at stellar masses greater than $10^{10} M_\odot$, but at lower stellar masses the models predict a weaker bias for gas-deficient galaxies than is actually seen.

As seen in the bottom panels of Fig. 8, the change in bias factor between very gas-rich galaxies and galaxies with typical gas fractions appears to be weaker rather than stronger at low stellar masses. The most gas-rich galaxies with stellar masses greater than $10^{10} M_\odot$ are more weakly clustered than galaxies with typical gas fractions, indicating that they occupy lower mass dark matter haloes on average. At stellar masses below $10^{10} M_\odot$, there is no antibias of gas-rich galaxies seen in the data. However, the models do predict clear antibias effects.

Figure 8. In the upper panels, we plot the change in bias factor from the 10th to the 50th percentile in $\log_{10}(M_{\text{HI}}/M_\ast)$ as a function of projected physical scale, for different stellar mass intervals as indicated. The black line shows the result from the SDSS/DR7. The errors estimated from the mock catalogues are shown as shaded regions. The red dot–dashed line and the blue solid line show results from the Fu et al. (2010) and Guo et al. (2011) models after convolution with errors. The lower panels show the change in bias factor from the 50th to the 100th percentile.
One might question whether ranking galaxies by H\textsc{i} mass fraction is sufficient to characterize whether a galaxy is classified as gas-rich or gas-deficient. As discussed in Section 3, galaxies of fixed stellar mass and colour have higher H\textsc{i} mass fractions if they have larger sizes (i.e. lower stellar surface mass densities). One way to understand this is to appeal to standard disc formation models (e.g. F10; Kauffmann 1996; Mo, Mao & White 1998). In these models, the spin parameter of the dark matter halo determines the contraction factor of the infalling gas. Larger discs in a dark matter halo of fixed mass will have higher H\textsc{i} mass fractions because gas surface densities are low and gas consumption times are long. In this case, it would make more sense to define galaxies as gas-rich or gas-deficient by comparing their H\textsc{i} mass fractions to other galaxies of the same mass and size.

One might also consider an even more stringent constraint that H\textsc{i}-rich/H\textsc{i}-deficient galaxies be classified as those objects with higher/lower than average H\textsc{i} content given their stellar mass, size and star formation rate. This might indicate that the galaxy has experienced a recent gas accretion episode and that the global star formation has not yet had a chance to respond to the extra fuel supply. In our scheme of using photometric quantities to predict H\textsc{i} content, the H\textsc{i}-rich systems would correspond to those galaxies with bluer than average outer discs. Recall that the H\textsc{i} content in gas-poor regime is currently calibrated using only stellar surface density and colour; we therefore do not delve into the opposite regime, where gas has been recently removed from a galaxy.

In Figs 9 and 10 we investigate clustering trends using these alternative definitions. For Fig. 9, we rank galaxies as a function of their deviation from the average H\textsc{i} mass fraction of all galaxies of the same stellar mass (M*) and stellar surface mass density (\(\mu_\ast\)). As seen from equation (1), this deviation depends on both the NUV - r colour of the galaxy and its g - i colour gradient. For Fig. 10, we rank galaxies as a function of their deviation from the average H\textsc{i} mass fraction of all galaxies of the same M*, \(\mu_\ast\) and NUV - r. This then depends only on the g - i colour gradient of the galaxy.

Interestingly, the top panels of Fig. 9 show that when gas deficiency is expressed relative to galaxies of the same mass and size, the change in bias on scales between a few hundred kpc and 1 Mpc becomes much more pronounced. The change in bias factor for the lowest mass galaxies now reaches values near \(\sim 3\) and even massive gas-deficient galaxies are now significantly biased with respect to their counterparts with ‘normal’ gas fractions.

The F10 model provides predictions of the radial profiles of the gas and the stars in galaxies. It is thus possible to look at gas deficiency at fixed mass and stellar surface density in the model. Results are plotted as red curves in Fig. 9. We see that the model disagrees very strongly with the observations. In the model, bias effects become weaker rather than stronger when gas deficiency is defined with respect to galaxies of the same mass and size. These results would appear to suggest that in the real Universe, gas removal processes depend on the size/density of the galaxy itself. This is not the case in the models, where satellite galaxies become gas-poor only because their supply of infalling gas has been cut off. Thus the data suggest that processes such as ram-pressure stripping, which depend on the density of the ISM, may play an important part in explaining the observed trends.

In contrast to what is seen for gas-deficient galaxies, if gas richness is normalized with respect to galaxies of the same stellar mass and size, the bias trends remain much the same and the F10 model predictions still fit reasonably well for galaxies more massive than \(10^{10}\, M_\odot\). This suggests that gas accretion processes are being modelled quite successfully at high stellar masses.

Fig. 10 shows that the bias effects for gas-rich galaxies are still roughly of the same strength as in the previous figure, when gas...
richness is expressed relative to galaxies of the same mass, size and global NUV - r colour. This means that the clustering does depend strongly on the colour-gradient term for these objects. Galaxies with bluer than average outer colours are clearly located in lower density environments compared to galaxies where there is no evidence for younger than average stellar populations in the outer disc.

6 SUMMARY AND DISCUSSION

We introduce a new photometric estimator for estimating the H I mass fraction \( M_{\text{HI}}/M_\star \) in local galaxies. The estimator is calibrated with a sample of 293 galaxies with \( M_\star > 10^{10} \, M_\odot \) in the redshift range 0.025 < \( z < 0.05 \), which are detected in HI emission line by the GASS survey. The estimator is a linear combination of four parameters: stellar mass \( M_\star \), stellar surface mass density \( \mu_* \), near-UV-to-optical colour \( NUV - r \) and the gradient in \( g - i \) colour \( \Delta g - \Delta i \). We demonstrate that this estimator provides unbiased \( M_{\text{HI}}/M_\star \) estimates for H I-rich galaxies.

We then apply this estimator to a sample of \( \sim 24,000 \) galaxies from the SDSS/DR7 that lie in the same redshift range. We analyse the clustering of these galaxies as a function of stellar mass and as a function of H I mass fraction \( M_{\text{HI}}/M_\star \) and we compare the results with predictions from two recent semi-analytic models of galaxy formation by F10 and G11. Our results may be summarized as follows.

(i) Clustering depends strongly on H I mass fraction at fixed \( M_\star \). Galaxies with large values of \( M_{\text{HI}}/M_\star \) are more weakly clustered. The total change with H I mass fraction in clustering strength is largest for low-mass galaxies.

(ii) At fixed \( M_\star \), the clustering dependence on H I mass fraction is strongest on scales of a few hundred kpc. On large scales (>1 Mpc), clustering depends weakly on H I mass fraction. This suggests that the H I content of a galaxy of fixed stellar mass depends on location within its dark matter halo.

(iii) After the uncertainty in the H I mass fraction estimator is taken into account, the observed dependence of clustering on H I mass fraction is well reproduced by the models for galaxies more massive than \( 10^{10} \, M_\odot \). Significant discrepancies remain at lower stellar masses.

In the next part of the paper, we extend the analysis by studying the clustering of H I-deficient and H I-rich galaxies defined in two ways: (1) with respect to the average H I content of galaxies of the same stellar mass and size, and (2) with respect to the average H I content of galaxies of the same stellar mass, size and NUV - r colour. These definitions are motivated by the following considerations. First, models in which discs form from gas that cools and condenses in dark matter haloes, while conserving angular momentum, predict that the gas fractions of equilibrium discs depend on both their mass and their size. Secondly, the majority of nearby spiral galaxies are observed to lie on a relatively tight plane linking H I gas mass fraction with stellar mass, stellar surface density and NUV - r colour (C10). It is natural to suppose that galaxies that have undergone a recent gas accretion event would be displaced to higher H I mass fractions with respect to this plane. Conversely, galaxies that have been stripped of their gas would be displaced to lower values of \( M_{\text{HI}}/M_\star \).

The main results of our analysis of H I-rich and H I-deficient galaxies can be summarized as follows.

(i) When H I deficiency is defined with respect to galaxies of the same stellar mass and size, the bias of H I-deficient galaxies relative to normal galaxies is larger than obtained if the H I deficiency is defined with respect to galaxies of the same mass. The same effect is not reproduced in the semi-analytic model of F10.

(ii) When H I richness is expressed with respect to galaxies of the same mass and size (as well as with respect to galaxies of the same mass, size and colour), H I-rich galaxies more massive than \( 10^{10} \, M_\odot \) are observed to be antibiased with respect to their counterparts with normal H I content. The same is not true at lower stellar masses.

We have proposed that the disagreement between the observations and the models might be resolved, if processes such as ram-pressure stripping, which depend on the density of the ISM, are included in the models. We note that the lowest mass galaxies have the lowest densities and are thus the most likely to be affected by ram pressure. In order to test this hypothesis in more detail, we plan to look at the behaviour of gas deficiency as a function of cluster-centric radius in samples of nearby groups and clusters.

We also stress that next generation wide-field H I surveys such as the Australian Square Kilometre Array Pathfinder H I All-sky Survey (WALLABY) and surveys carried out by the Apertif receiver array on the Westerbork Synthesis Radio Telescope will measure HI masses and sizes for samples of tens to hundreds of thousands of galaxies at redshifts of around 0.1. These surveys will make it possible to investigate clustering as a function of the true H I content of a galaxy. It will be interesting to investigate the degree to which the results conform with our current analysis of 'pseudo'-H I content. Discrepancies may reveal additional physics that we have not yet considered. In the meantime, the construction of models that can reproduce the gas properties of galaxies as well as possible is...
an important step towards building mock surveys that can be safely used for making predictions in support of these surveys.

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