DEPENDENCE OF THE FREQUENCY OF THE KILOHERTZ QUASI-PERIODIC OSCILLATIONS ON X-RAY COUNT RATE AND COLORS IN 4U 1608–52

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ABSTRACT

We present new results based on observations carried out with the Rossi X-ray Timing Explorer during the decay of an outburst of the low-mass X-ray binary (LMXB) and atoll source 4U 1608–52. Our results appear to resolve, at least in 4U 1608–52, one of the long-standing issues about the phenomenology of the kilohertz quasi-periodic oscillations (kHz QPOs), namely the lack of a unique relation between the frequency of the kHz QPOs and the X-ray flux. We show that despite its complex dependence on the X-ray flux, the frequency of the kHz QPOs is monotonically related to the position of the source in the color-color diagram. Our findings strengthen the idea that, as in the case of Z sources, in the atoll sources the X-ray flux is not a good indicator of and that the observed changes in the frequency of the kHz QPOs in LMXBs are driven by changes in M. These results raise some concern about the recently reported detection of the orbital frequency at the innermost stable orbit in 4U 1820–30.

Subject headings: accretion, accretion disks — stars: individual (4U 1608–52) — stars: neutron — X-rays: stars

1. INTRODUCTION

Nearly 3 years have elapsed since the first kilohertz quasi-periodic oscillations (kHz QPOs) were discovered with the Rossi X-ray Timing Explorer (RXTE) in the X-ray flux of Scorpius X-1 (van der Klis et al. 1996) and 4U 1728–34 (Strohmayer, Zhang, & Swank 1996). In the meantime, kHz QPOs have been observed in the persistent flux of 18 low-mass X-ray binaries (see van der Klis 1998 for a review), both in the so-called Z sources and in the atoll sources (Hasinger & van der Klis 1989, hereafter HK89). Except in Aql X-1, which showed a single kHz peak in its power spectrum (Zhang et al. 1998a), in the other 17 sources two simultaneous kHz QPO peaks have been observed. In some sources a third, nearly coherent, QPO peak has been detected during type I X-ray bursts, with a frequency that was consistent with being equal to 1 or 2 times the frequency separation of the kHz QPOs observed in the persistent flux (Strohmayer, Swank, & Zhang 1998). It was suggested that a beat frequency mechanism is responsible for this commensurability in the kHz QPO frequencies (Strohmayer et al. 1996; Miller, Lamb, & Psaltis 1998), but this interpretation is not without problems (van der Klis et al. 1997; Méndez et al. 1998c; Méndez, van der Klis, & van Paradijs 1998b).

The dependence of the kHz QPO frequencies on X-ray luminosity, which is usually assumed to be a measure of the mass accretion rate, is complex. While in a given source there is a good correlation between frequency and luminosity on a timescale of hours, sources that span nearly 3 orders of magnitude in luminosity, such as Sco X-1 and 4U 0614+09, show kHz QPOs that cover the same range of frequencies (van der Klis 1997). It is as if the frequency of the kHz QPO depends on the difference between the instantaneous and average luminosity in each source rather than on the luminosity itself.

A similar effect is seen between observations of the same source at different epochs. On timescales of hours, frequency and X-ray flux are well correlated, but between different epochs the source covers the same range of frequencies even if the average flux is different by 40% or more (e.g., Aql X-1; Zhang et al. 1998a).

In this Letter, we present new results that appear to resolve the latter of these two issues, at least in 4U 1608–52. We show that while on timescales longer than ~1 day the frequency of the kHz QPOs is not well correlated to the X-ray flux, it is very well correlated to the position of the source in the color-color diagram. From this result, we conclude that the observed changes in the frequency of the kHz QPOs in 4U 1608–52 are driven by changes in the mass accretion rate and that the lack of correlation between QPO frequency and X-ray count rate occurs because, as in the case of the Z sources, in atoll sources there is no one-to-one relationship between the observed X-ray flux and the mass accretion rate.

2. OBSERVATIONS AND RESULTS

All of the observations presented here were obtained using the proportional counter array (PCA) on board RXTE. We include the data of the decay of the 1996 outburst (Berger et al. 1996; Méndez et al. 1998a) and of the 1998 outburst (Méndez et al. 1998c; Fig. 1 there shows a light curve of the 1998 outburst.) The observations as well as the modes used to record the data are described in Méndez et al. (1998c). We also include here a recent public target-of-opportunity RXTE/PCA observation of ~7.3 ks performed on 1998 August 6, 03:13 UTC. The observing modes for this last observation were similar to those used by Méndez et al. (1998c) after 1998 March 27.

We calculated count rates in five energy bands (2.0–3.5, 3.5–6.4, 6.4–9.7, 9.7–16.0, and 2.0–16.0 keV), taking into account the gain changes applied to the PCA in 1996 March and April. In a few of the observations, one or two of the five detectors of the PCA were switched off; we only used the three detectors that were always on to calculate these count rates. We subtracted the background contribution in each band using
the standard PCA background model version 2.0c and normalized the count rates to five detectors.

In Figure 1 we show a color-color diagram of 4U 1608–52. This is the first time that 4U 1608–52 is observed to move across all of the branches of the atoll, and this constitutes one of the best examples of the color-color diagram of an atoll source. Based on this diagram, we conclude that in 1996, as the source count rate decreased, 4U 1608–52 gradually moved from the lower part of the banana to the island state. In 1998, at the peak of the outburst, 4U 1608–52 was in the upper part of the banana and gradually moved down to the lower part of the banana and the island state as the count rate decreased. In general, the count rate is observed to increase along the track from the island state to the banana branch, but the relation between count rate and either of the two colors is much less clean than that between colors.

The high time resolution data confirm this preliminary state classification. We divided the 2–60 keV data into segments of 256 and 512 s and calculated power spectra up to a Nyquist frequency of 2048 Hz, normalized to fractional rms squared per hertz. The characteristics of the ≈100 Hz part of these power spectra changed in correlation with the position of the source in the color-color diagram. When 4U 1608–52 was in the upper and lower parts of the banana, the power spectra fitted a power law below ~1 Hz (the very low frequency noise, VLFN) and an exponentially cutoff power law above ~1 Hz (the high-frequency noise, HFN; see HK89). As 4U 1608–52 moved from the upper to the lower parts of the banana, the fractional amplitude of the VLFN (0.001–1 Hz) decreased from ~6% to ~2% rms, while the fractional amplitude of the HFN (1–100 Hz) increased from ~1% to ~5% rms. In the island state, the VLFN disappeared completely (the 95% confidence upper limits were ~1% rms), and the amplitude of the HFN increased further to ~10%–17% rms. (A more detailed analysis of the low-frequency part of the power spectra will be presented elsewhere.)

The ≈100 Hz part of the power spectra also changed in correlation with the position of the source in the color-color diagram: we only observed kHz QPOs when 4U 1608–52 was at certain positions of the color-color diagram (see Fig. 1). For those segments in which we observed QPOs, the 2–60 keV fractional amplitudes of the lower frequency and higher frequency, hereafter the lower and upper QPO, varied from 5.3% to 9.1% rms and from 3.3% to 8.8% rms, respectively. We did not detect kHz QPOs in the upper part of the banana, with 95% confidence upper limits of 0.8%–4.6% rms depending on the source count rate and the assumed width of the QPO or in the extreme island state (upper right-hand corner of the color-color diagram), with 95% confidence upper limits of 3.5%–10% rms. While these upper limits strongly suggest the absence in the upper banana of kHz QPOs as strong as those observed in the lower banana and the island states, we cannot rule out the presence of such kHz QPOs in the extreme island state, when the count rates were lowest.

To further characterize the dependence of the kHz QPOs on other source parameters, we selected only those data for which we detected two simultaneous kHz QPO peaks in the power spectrum (see Méndez et al. 1998c), so that the identification of the observed kHz peaks is unambiguous. We divided the data into segments of 64 s and produced a power spectrum for each segment extending from 1/64 to 2048 Hz. In Figure 2 we show the dependence of the frequency of the lower QPO $f_{low}$.
as a function of count rate for 4U 1608–52. This figure shows several branches, which reflect the relation between \( n_{\text{low}} \) and count rate during individual observations that span from \( \sim 0.5 \) to \( \sim 8 \) hr. The only exception is the branch at the lowest count rate, where two different observations taken 8 days apart overlap. This figure clearly shows that while on timescales of a few hours or less \( n_{\text{low}} \) is well correlated to count rate, on timescales greater than \( \sim 1 \) day the relation is complex and \( n_{\text{low}} \) is not uniquely determined by the count rate. We obtained the same result using the 2–60 keV source flux instead of the count rate.

There is a much better correlation between \( n_{\text{low}} \) and the position of the source on the color-color diagram. In Figure 3 we show \( n_{\text{low}} \) as a function of hard color (see Fig. 1) for the same segments shown in Figure 2. The complexity seen in the frequency versus count rate diagram (Fig. 2) is reduced to a single track in the frequency versus hard color diagram. The frequency of the lower kHz QPO increases as the hard color decreases, i.e., as 4U 1608–52 moves from the island state to the lower banana, and keeps increasing at the turn of the lower banana, where the hard color reaches its lowest value of \( \sim 0.40 \).

The shape of the track in Figure 3 suggests that the hard color may not be sensitive to changes of state when the source moves into the banana in the color-color diagram. To further investigate this, we applied the \( S_{a} \) parameterization (e.g., Wijnands et al. 1997b), which we call \( S_{a} \) for atoll sources, to the color-color diagram. In this phenomenological approach we approximated the shape of the color-color diagram with a spline (Fig. 1), and we used the parameter \( S_{a} \) to measure positions along this spline. We set \( S_{a} \) to 1 at (2.67, 0.77) and to 2 at (2.19, 0.42), as indicated in Figure 1. We measured the position on the color-color diagram of each of the 64 s segments. We then grouped these segments into 36 sets, such that within each set \( n_{\text{low}} \) did not vary by more than \( \sim 25 \) Hz and the source count rate did not vary by more than 10%. (This is to avoid mixing segments that come from different branches in Fig. 2.) For each set we calculated the average frequency of the lower and upper kHz QPO peak and of \( S_{a} \). In Figure 4 we plot the frequencies of both kHz QPOs as a function of \( S_{a} \). The error bars are the standard deviation of each selection. We also include in this figure several extra sets (open squares) for which we only detected one of the kHz QPOs, which therefore we could not a priori identify as the upper or lower peak. As expected, \( S_{a} \) is more sensitive than the hard color to changes of state when the source moves from the island state into the turn of the banana: \( S_{a} \) increases from \( \sim 2.04 \) to \( \sim 2.15 \) there, while the hard color saturates at \( \sim 0.40 \). Figure 4 confirms that the frequency of the kHz QPOs is strongly correlated to the position of the source in the color-color diagram. It also shows that the steep rise of kHz QPO frequency at the turn of the lower banana is not just due to a lack of sensitivity of the hard color to state changes in this part of the color-color diagram but that it occurs as a function of the position in the color-color diagram as well.

3. DISCUSSION

Previous observations of kHz QPOs in atoll sources have resulted in a confusing picture concerning the dependence of QPO frequency on mass accretion rate. Cases have been reported in which, in a given source, QPO frequency showed a good correlation to count rate or spectral hardness (Ford, van der Klis, & Kaaret 1998a; Wijnands et al. 1997c; Wijnands et al. 1998; Ford et al. 1998b; Zhang et al. 1998b; Wijnands & van der Klis 1997). In other cases, a correlation with count rate or flux was conspicuously lacking (Ford et al. 1997; Méndez et al. 1998a; Zhang et al. 1998a).

Our observations show for the first time that a total lack of correlation between frequency and count rate on timescales longer than a day (Fig. 2) can coexist with a very good correlation between frequency and position in the X-ray color-color diagram (Fig. 3). The frequency increases with \( S_{a} \) as the source moves from the island to the banana. Only on timescales of hours does the QPO frequency appear to also correlate well.
with count rate. The presence of the QPOs also correlates well with the position in the color-color diagram: the QPOs are only detected in the lower banana and the moderate island states and disappear both in the upper banana and in the extreme island states (Fig. 1).

In atoll sources, $M$ is thought to increase monotonically with $S_\lambda$ along the track in the color-color diagram, from the island to the upper banana (HK89), whereas X-ray count rate tracks $M$ much less well (van der Klis et al. 1990; van der Klis 1994; Prins & van der Klis 1997). The properties of the power spectra below 100 Hz depend monotonically on inferred $M$ (HK89). Our result that the frequency of the kHz QPO is well correlated to $S_\lambda$ but not to X-ray count rate implies that in 4U 1608–52 the kHz QPO frequency also depends monotonically on inferred $M$. By extension, this conclusion applies to each atoll source; in Z sources, similar conclusions were previously also reached (e.g., Wijnands et al. 1997a). Our analysis, however, sheds no light on the question of why sources with very different inferred $M$’s (e.g., 4U 0614+09 and Sco X-1) can have kHz QPOs in the same range of frequencies.

Further indications for this interpretation come from the simultaneous analysis of the low- and high-frequency parts of the power spectra of other atoll sources. In 4U 1728–34, the kHz QPO frequencies were recently found to be very well correlated to several properties of the power spectra below 100 Hz (Ford & van der Klis 1998), and Psaltis, Belloni, & van der Klis (1998) obtained a similar result for a number of other atoll (and Z) sources. In all of these sources, not only the position in the color-color diagram and the various low-frequency power spectral parameters, but also the frequencies of the kHz QPOs are all well correlated with each other. This indicates that the single parameter, inferred to be $M$, which governs all of the former properties also governs the frequency of the kHz QPO.

X-ray intensity is the exception: it can vary more or less independently from the other parameters. In 4U 1608–52, it can change by a factor of $\sim 4$ (see Fig. 2), while the other parameters do not vary significantly. As inferred, this constancy of the other parameters means that $M$ does not change, then this indicates that strongly variable beaming of the X-ray flux or large-scale redistribution of some of the radiation over unobserved energy ranges is occurring in order to change the flux by the observed factors, without any appreciable changes in the X-ray spectrum. We may need to scrutinize more closely the concept of $M$ in order to solve this dilemma. For example, perhaps the $M$ governing all of the other parameters is the $M$ through the inner accretion disk, whereas matter also flows onto the neutron star in a more radial inflow or away from it in a jet.

In 4U 1608–52, we observe no evidence for a saturation of the frequency of the kHz QPOs at a constant maximum value as $M$ increases, different from what Zhang et al. (1998b) inferred for 4U 1820–30. They presented data in which the kHz QPO frequencies increase with count rate up to a threshold level, above which the frequencies remain approximately constant while the count rate keeps increasing. Interpreting count rate as a measure for $M$, they argue that this is evidence for the inner edge of the disk reaching the general relativistic innermost stable orbit. However, we have shown here that, at least in 4U 1608–52, count rate is not a good measure for $M$.

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