A Three-Dimensional Random Walk Algorithm for Estimating the Chloride Diffusivity of Concrete

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Abstract: The chloride diffusivity of concrete is an important parameter for assessing the long-term durability of coastal concrete structures. The purpose of this paper is to present a three-dimensional random walk algorithm (RWA) for estimating the chloride diffusivity of concrete. By analyzing the size distribution of aggregates, the equivalent interfacial transition zone (ITZ) thickness is derived in an analytical manner. Each aggregate is combined with the surrounding ITZ to construct an equivalent aggregate model (EAM) and the chloride diffusivity is formulated. It is found that the equivalent ITZ thickness decreases with the increase of practical ITZ thickness and aggregate volume fraction. The aggregate gradation influences the equivalent ITZ thickness to a certain extent. The relative chloride diffusivity of the equivalent aggregate is almost directly and inversely proportional to the equivalent ITZ thickness and the aggregate radius, respectively. The numerical results show that, when the EAM is adopted, the computational time is greatly reduced. With the EAM, concrete can be modeled as a two-phase material and the chloride diffusivity is estimated by applying the RWA. It is shown that, with the increase of mean square displacement and number of Brownian particles, the average chloride diffusivity of concrete approaches a stable value. Finally, through comparison with experimental data, the validation of the RWA is preliminarily verified.

Keywords: concrete; chloride diffusivity; EAM; RWA; spherical aggregate

1. Introduction

Reinforced concrete structures built in a chloride-laden environment often suffer from corrosion of reinforcement caused by penetration of chloride ions [1,2]. Many such engineering cases have been reported throughout the world [3–5]. Since the long-term performance of concrete infrastructure is greatly dependent on the diffusion rate of chloride ions in concrete, the chloride diffusivity plays a key part in designing and assessing coastal concrete structures [6,7].

In the past twenty years, a considerable number of experimental studies and theoretical analyses have been performed with a particular focus on the transport properties of concrete and on the key influential factors. By taking the water/cement ratio and the sand content as the primary control variables, three types of mortars were examined by a diffusion test and a migration test [6]. The test results showed that addition of sand into cement paste modifies the microstructure and chloride diffusivity of mortar. The tortuosity caused by sand particles exerts a more significant effect on the diffusion of chloride ions in mortar than the interfacial transition zone (ITZ). Using an accelerated
method, Yang and Su quantified the effects of dilution, tortuosity, and ITZ on the chloride diffusivity of concrete based on the Bruggeman theory [8]. Caré [9] adopted a non-steady-state diffusion test to further analyze the relative importance between two competing factors: ITZ and tortuosity. In the theoretical aspect, a Padé approximation was proposed to estimate the chloride diffusivity of mortar [10]. An analytical solution was also presented for the chloride diffusivity of concrete with low aggregate volume fractions [11]. Caré and Herve proposed a composite sphere model for the chloride diffusivity of mortar [12]. The effects of various factors on the chloride diffusivity of concrete were comprehensively evaluated by a transfer matrix method with an inhomogeneous ITZ model [13]. It was also shown that aggregate shape influences the chloride diffusivity of concrete to a certain extent [14]. Besides a few analytical methods [15,16], the unit cell model is widely used, in which the volume ratios among various phase constituents are identical. In practical concrete, however, the ITZ thickness seems to be the same no matter how large the aggregate is [17]. As a consequence, a larger aggregate is of a smaller ITZ volume ratio in a unit cell. Thus, the prediction accuracy of the idealized unit cell with constant volume ratios will be affected to a certain extent. Although numerical methods can avoid this defect by simulating the heterogeneous structure of concrete, it is very time-consuming [18,19]. Therefore, it is still essential to develop a numerical method, with which the chloride diffusivity of concrete can be evaluated more reasonably.

This paper is aimed at developing a three-dimensional random walk algorithm (RWA) for estimating the chloride diffusivity of concrete. In the RWA, an equivalent aggregate model (EAM) is constructed to reduce the computational cost. Finally, the computational accuracy of the RWA is assessed through comparison with experimental results.

2. Simulation of Concrete Mesostructure

In what follows, only the steady-state diffusion of chloride ions in saturated concrete is considered. It is appreciated that transport of chloride ions in the pore solution is affected, to a certain extent, by the interactions among ions and the electrical double layer [20–23]. In view of the complexity of the two effects, they are neglected as a first step. To perform the random walk simulation, the concrete mesostructure needs to be reproduced as realistically as possible. For this purpose, a cubic element with a side length of a is selected. Aggregates are modeled as spherical and distributed at random within the element [24]. If the aggregate is divided into N grades \([R_j, R_{j+1}]\) (\(j = 1, 2, \ldots, N\)) and the volume percentage passing the sieve with radius \(R_j\) is \(P_{vj}\), the approximate probability density function \(p_v(R)\) with respect to the volume of aggregates is given by

\[
p_v(R) = \sum_{j=1}^{N} \frac{P_{vj+1} - P_{vj}}{R_{j+1} - R_j} [H(R - R_j) - H(R - R_{j+1})], \tag{1}
\]

where \(H(x)\) is the Heaviside step function and defined as

\[
H(x) = \begin{cases} 
1, & x > 0 \\
0.5, & x = 0 \\
0, & x < 0
\end{cases} \tag{2}
\]

The corresponding cumulative distribution function \(P_v(R)\) is given by

\[
P_v(R) = \int_{R_1}^{R} p_v(x)dx = P_{v,k} + \sum_{k=1}^{N} \frac{P_{vk+1} - P_{vk}}{R_{k+1} - R_k} (R - R_k), \tag{3}
\]
where \( R_k \leq R \leq R_{k+1} \) and \( k = 1, 2, \ldots, N \). The number of aggregates per unit volume of aggregate can be formulated as

\[
N_v = \frac{3}{4\pi} \int_{R_1}^{R_{N+1}} \frac{p_v(x)}{x^3} \, dx. \tag{4}
\]

By substituting Equation (1) into Equation (4), one has

\[
N_v = \frac{3}{8\pi} \sum_{j=1}^{N} \frac{(R_j + R_{j+1})(P_{v,j+1} - P_{v,j})}{R_j^2 R_{j+1}^2}. \tag{5}
\]

It is easily shown that the probability density function \( p_n(R) \) with respect to the number of aggregates is given by

\[
p_n(R) = \frac{3}{4\pi} \cdot \frac{p_v(R)}{N_v R^3}. \tag{6}
\]

It follows from Equations (1) and (6) that

\[
p_n(R) = \frac{3}{4\pi} \sum_{j=1}^{N} \frac{(P_{v,j+1} - P_{v,j})}{N_v (R_{j+1} - R_j) R^3} \left[ H(R - R_j) - H(R - R_{j+1}) \right]. \tag{7}
\]

The corresponding cumulative distribution function \( P_n(R) \) is given by

\[
P_n(R) = \int_{R_1}^{R} p_n(x) \, dx
\]

\[
= \frac{3}{8\pi} \sum_{j=1}^{k-1} \frac{(R_j + R_{j+1})(P_{v,j+1} - P_{v,j})}{N_v R_j^2 R_{j+1}^2} + \frac{(R^2 - R_k^2)(P_{v,k+1} - P_{v,k})}{N_v (R_{k+1} - R_k) R_k^2 R^2}, \tag{8}
\]

where \( R_k \leq R \leq R_{k+1} \) and \( k = 1, 2, \ldots, N \).

Knowing \( P_n(R) \) and the aggregate volume fraction \( f_a \), the aggregates to be distributed within the cubic element can be generated [24]. These aggregates are placed into the element from largest to smallest and no overlap is permitted between them. During the placement process, periodic boundary conditions are imposed to eliminate artificial wall effects. Finally, the reconstruction of concrete mesostructure is completed once each aggregate is surrounded with an ITZ of thickness \( h \).

3. Equivalent Aggregate Model and ITZ Thickness

It has been shown that there are two opposite effects on the diffusion of chloride ions in concrete. The dilution and tortuosity induced by aggregates decelerate the movement of chloride ions. On the other hand, the ITZ has a larger water/cement ratio (w/c) and higher porosity compared with the bulk cement paste. Furthermore, for a sufficiently high aggregate volume fraction, the ITZ percolates throughout the concrete specimen [25]. Thus, the ITZ and percolation effects accelerate the movement of chloride ions. Since the ITZ reduces the water/cement ratios of neighboring zones and the effect is internally balanced, its net effect on the movement of chloride ions is actually small [13]. To reduce the computational cost [26], an EAM is constructed as follows.

Since the microstructure of ITZ is different from that of bulk cement paste, it is more reasonable to model the ITZ as a distinct phase. Compared with the aggregate size, the ITZ thickness is usually much smaller [27]. As a result, when a Brownian particle walks near an ITZ, it will spend a great deal of time to walk even a short distance [18]. Therefore, the random walk simulation cannot efficiently be performed on the mesostructure of concrete.

To overcome this difficulty, an EAM is constructed by combining each aggregate with the surrounding ITZ, as shown in Figure 1. In practical concrete, when the surface-to-surface distance between two aggregates is smaller than twice ITZ thickness, their ITZ layers will overlap each other.
As a result, the ITZ thickness $h_{eq}$ in the EAM is smaller than the practical thickness $h$. To determine the equivalent ITZ thickness, the $k$th moment of area $<R^k>$ of $p_n(x)$ about the origin is defined as [28]

$$<R^k> = \int_{R_1}^{R_{N+1}} x^k p_n(x) dx. \tag{9}$$

It is evident that $<R>$, $4\pi <R^2>$, and $4\pi <R^3>/3$ represent the average radius, surface area, and volume of spherical aggregates, respectively. By substituting Equation (7) into Equation (9), one has

$$<R^k> = \begin{cases} \frac{3}{4\pi} \sum_{j=1}^{N} \frac{(P_{R_{j+1}} - P_{R_j}) \ln(R_{j+1}/R_j)}{N_t(R_j - R_j)}, & \text{for } k = 2 \\ \frac{2}{4\pi} \sum_{j=1}^{N} (P_{R_{j+1}} - P_{R_j})(R_{j+1}^k - R_j^k), & \text{for other cases} \end{cases}. \tag{10}$$

![Figure 1.](image)

**Figure 1.** (a) Three-phase concrete; (b) equivalent aggregate; and (c) two-phase concrete.

An algorithm was presented for evaluating the ITZ volume fraction $f_i$ based on the statistical geometry of composites [11,29]. According to this algorithm, $f_i$ is given by

$$f_i = (1 - f_a)[1 - \exp(-t_1 h - t_2 h^2 - t_3 h^3)] \tag{11}$$

where the coefficients $t_1$, $t_2$, and $t_3$ are expressed in terms of $f_a$ and $<R^k>$ as [29]

$$t_1 = \frac{3f_a<R^2>}{(1 - f_a)<R^3>} \tag{12a}$$

$$t_2 = \frac{3f_a<R>}{(1 - f_a)<R^3>} + \frac{9f_a^2<R^2>^2}{2(1 - f_a)^2<R^3>^2} \tag{12b}$$

$$t_3 = \frac{f_a}{(1 - f_a)<R^3>} + \frac{3f_a^2<R>^2}{(1 - f_a)^2<R^3>^2} + \frac{w f_a^3<R^3>^3}{(1 - f_a)^3<R^3>^3} \tag{12c}$$

with $w$ being 0, 2, or 3. It was found that the effect of $w$ on the computational accuracy is negligibly small but best results will be achieved for $w = 0$ [11]. Thus, $w$ is set to be zero in this paper. Since the ITZ volume fraction in the EAM should be equal to that in practical concrete, the equivalent ITZ thickness $h_{eq}$ satisfies the following equation

$$\frac{4\pi}{3} \int_{R_1}^{R_{N+1}} N_t f_a p_n(x)(3x^2 h_{eq} + 3xh_{eq}^2 + h_{eq}^3) dx = f_i. \tag{13}$$
By substituting Equation (9) into Equation (13), one has
\[ N_f f_a (h_{eq}^3 + 3(R)h_{eq}^2 + 3(R^2)h_{eq}) = \frac{3}{4\pi} f_i, \]  
(14)

Solving Equation (14) for \( h_{eq} \), one has
\[ h_{eq} = 3\sqrt{-q^2 + \sqrt{q^4 + p^3}} \]
(15)

where \( p \) and \( q \) are equal to
\[ p = 3(\langle R^2 \rangle - \langle R \rangle^2) \]  
(16a)
\[ q = 2(\langle R \rangle^3 - 3(\langle R \rangle^2 \langle R \rangle) - \frac{3}{4\pi} \cdot \frac{f_i}{N_f f_a}. \]  
(16b)

It will be seen below that \( h_{eq} \) is closely related to the ITZ volume fraction in each equivalent aggregate and therefore the chloride diffusivity of concrete. It is interesting to investigate the effects of \( f_a \), \( h \), and aggregate gradation on \( h_{eq} \), as seen from Equations (11), (15), and (16). To evaluate these influential factors in a quantitative manner, the Fuller gradation is first adopted with sizes from 0.15 to 16 mm. Thus, \( h_{eq}/h \) is plotted against \( f_a \) in Figure 2 [24,30], which demonstrates that, for a small value of \( f_a \), \( h_{eq}/h \) approaches unit and therefore ITZs seldom overlap. As \( f_a \) increases, more and more ITZs overlap and \( h_{eq}/h \) decreases. For a given \( f_a \), \( h_{eq}/h \) decreases with the increase of \( h \). This is due to the fact that a larger \( h \) results in more overlaps of ITZs. For a given \( f_a \) at 0.6, 0.7, and 0.8, \( h_{eq}/h \) at \( h = 0.05 \) mm is smaller than that at \( h = 0.01 \) mm by 5.76%, 10.1%, and 20.1%, respectively.

Second, the effect of aggregate gradation on \( h_{eq}/h \) is analyzed. For this purpose, two typical gradations, Fuller and equal volume fraction (EVF) [24,30], are considered and \( h = 0.03 \) mm. Thus, \( h_{eq}/h \) is plotted against \( f_a \) in Figure 3, indicating that, for a given \( f_a \) at 0.6, 0.7, and 0.8, \( h_{eq}/h \) for
concrete with the EVF gradation is smaller than that with the Fuller one by 5.29%, 9.82%, and 20.1%, respectively. This attributes the fact that more small aggregates in the EVF gradation result in more overlaps of ITZs.

Figure 3. Effect of aggregate gradation on $h_{eq}/h$.

4. Chloride Diffusivity of Concrete

Duan et al. [31] proposed a completely explicit formula for the conductivity tensor of multi-phase media with various inclusions. From the formula, the chloride diffusivity of the equivalent aggregate shown in Figure 1 is given by

$$D_{ea} = D_i + \frac{V_a}{V_a + V_i} \left( \frac{D_a - D_i}{D_i + (D_a - D_i)/3} \right)^{-1} - \frac{1}{3} \right)^{-1}$$

where $D_a$ and $D_i$ are the chloride diffusivities of aggregate and ITZ, respectively, and the aggregate volume $V_a$ and the ITZ volume $V_i$ are equal to

$$V_a = \frac{4\pi}{3} R^3$$  \hspace{1cm} (18a)

$$V_i = \frac{4\pi}{3} [(R + h_{eq})^3 - R^3].$$  \hspace{1cm} (18b)

The chloride diffusivity of bulk cement paste is denoted by $D_{bcp}$. Since $D_a$ is much smaller than $D_i$ and $D_{bcp}$, $D_a$ is set to be zero. Thus, Equation (17) becomes

$$D_{ea} = \frac{2V_i}{3V_a + 2V_i} D_i.$$  \hspace{1cm} (19)
By substituting Equations (18a) and (18b) into Equation (19), one has

\[ D_{ea} = \frac{2[(R + h_{eq})^3 - R^3]}{2(R + h_{eq})^3 + R^3}D_i. \]  

(20)

It is seen from Equation (20) that, besides \( h_{eq} \), \( D_{ea}/D_i \) is also dependent on the aggregate size. If \( R_1 \) and \( R_{N+1} \) are set to be 0.075 and 8 mm, respectively, \( D_{ea}/D_i \) is plotted against \( R \) in Figure 4, indicating that, since \( h_{eq} \) is much smaller than \( R \), \( D_{ea}/D_i \) is almost directly and inversely proportional to \( h_{eq} \) and \( R \), respectively, as also seen from Equation (20).

Figure 4. Relationship between \( D_{ea}/D_i \) and \( R \) for different values of \( h_{eq} \).

With the EAM, concrete can be regarded as equivalent aggregates dispersed in a cement paste matrix. In this way, the RWA can directly be applied to the two-phase concrete [18,19].

In the beginning, a Brownian particle is randomly placed at a point, called the origin \( o \), in concrete, as shown in Figure 5. Then, it walks step by step until hitting the spherical surface \( \Gamma \) of radius \( R_0 \), i.e., the mean square displacement, centered at \( o \) for the first time. During the whole walk process, two cases should be considered separately. When the Brownian particle arrives at a point whose distance from any interfaces is larger than the prescribed value (=0.001 mm in this paper), a maximum sphere centered at the Brownian particle tangent to the interface is created. A random point is selected on the surface. If the spherical radius is \( r_i \) and the chloride diffusivity of the phase included in the sphere is \( D^{(i)} \), the mean time \( t(r_i) \) for the Brownian particle to jump to the random point is [18]

\[ t(r_i) = \frac{r_i^2}{6D^{(i)}}. \]  

(21)
It should be pointed out that Einstein and Smoluchowski derived an equation for Brownian movement in one dimension only under the assumption of Maxwell-Boltzmann distribution [32]. Later, the equation was extended to three dimensions. The theory of Brownian motion is usually developed for a system with no real boundaries. When the Einstein–Smoluchowski equation is directly applied to the diffusion of chloride ions in cementitious materials, the Brownian particle is restricted to three orthogonal directions [33]. When the Brownian particle encounters an impermeable solid phase, it is not permitted to step into the solid phase but the time spent is still counted for such an attempt [33]. There are two disadvantages of applying the Einstein–Smoluchowski equation to cementitious materials. First, this RWA simulates the detailed zigzag walk of the Brownian particle with small finite steps and is at least an order of magnitude slower than the Torquato and Kim one [34]. Second, when the Brownian particle comes near an interface between two permeable phases, it is difficult to compute the mean time and probability for it to cross the interface. To overcome these difficulties, Torquato and Kim [18] derived Equation (21) using the first passage time probability distribution. Based on Equation (21), the Brownian particle can walk directly to a random point on the surface of the maximum sphere, as stated above. Thus, there is no need to simulate the detailed zigzag walk of the Brownian particle with finite step sizes. Furthermore, when the Brownian particle is near an interface, they also formulated the mean time and probability for it to walk through the interface, as stated below.

When the Brownian particle arrived at a point x whose distance from an interface is smaller the prescribed value, the mean time can be evaluated as follows. If the projection of x on to the interface is denoted by \( x_0 \), a sphere of radius \( r \) is then created centered at \( x_0 \), as shown in Figure 6. The interface divides the sphere into two domains \( \Omega^{(1)} \) and \( \Omega^{(2)} \) with volumes \( V^{(1)} \) and \( V^{(2)} \) and chloride diffusivities.
D\(^{(1)}\) and D\(^{(2)}\) and the surface into two parts \(\Gamma^{(1)}\) and \(\Gamma^{(2)}\) with surface areas \(A^{(1)}\) and \(A^{(2)}\), respectively. The probabilities \(p_1\) and \(p_2\) for the Brownian particle to reach the surfaces \(\Gamma^{(1)}\) and \(\Gamma^{(2)}\) are equal to [18]

\[
\begin{align}
 p_1 &= \frac{A^{(1)}D^{(1)}}{A^{(1)}D^{(1)} + A^{(2)}D^{(2)}} \quad (22a) \\
 p_2 &= \frac{A^{(2)}D^{(2)}}{A^{(1)}D^{(1)} + A^{(2)}D^{(2)}} \quad (22b)
\end{align}
\]

![Figure 6. Brownian particle located near interface between phases 1 and 2.](image)

The mean time \(t(r_j)\) for the Brownian to jump to a random point on the spherical surface is formulated as

\[
t(r_j) = \frac{(V^{(1)} + V^{(2)})r_j^2}{6(V^{(1)}D^{(1)} + V^{(2)}D^{(2)})}. \quad (23)
\]

It follows from Equations (21) and (23) that the mean time \(t(R_0)\) is obtained as

\[
t(R_0) = \sum_i \frac{r_i^2}{6D^{(i)}} + \sum_j \frac{(V^{(1)} + V^{(2)})r_j^2}{6(V^{(1)}D^{(1)} + V^{(2)}D^{(2)})}. \quad (24)
\]

On the other hand, when concrete is homogenized, it becomes a uniform medium of chloride diffusivity \(D_{\text{con}}\), as shown in Figure 7. Thus, for a sphere of radius \(R_0\), the mean time \(t(R_0)\) for the Brownian particle to walk from the spherical center to a random point on the spherical surface is obtained as [18]

\[
t(R_0) = \frac{R_0^2}{6D_{\text{con}}}. \quad (25)
\]
On the other hand, when concrete is homogenized, it becomes a uniform medium of chloride diffusivity $D_{con}$, as shown in Figure 7. Thus, for a sphere of radius $R_0$, the mean time $t(R_0)$ for the Brownian particle to walk from the spherical center to a random point on the spherical surface is obtained as \[ t(R_0) = \frac{R_0^2}{\sum_i \frac{r_i^2}{D_i^0} + \sum_j \frac{(V^{(1)} + V^{(2)})}{\sum_i (D^{(1)} + D^{(2)})^2}}. \] (26)

Based on the developed algorithm, a computer program is written with FORTRAN language for reconstructing the mesostructure of concrete and for implementing the random walk of Brownian particles.

During the whole random walk process, the Brownian particle possibly walks outside the cubic element. In such a situation, periodic boundary conditions are imposed. When it walks across the face BCC$_1$B$_1$ shown in Figure 8a, the edge BC shown in Figure 8b, or the vertex C shown in Figure 8c, the line segment $yz$, which extends beyond the element boundaries, will be reflected into the element on the opposite face ADD$_1$A$_1$, the edge A$_1$D$_1$, or the vertex A$_1$, respectively, i.e., the line segment $y_1z_1$. When it walks across the other faces, edges, or vertexes, the procedure can be performed in a similar manner. To demonstrate the detailed walk process of a Brownian particle, a two-dimensional simulation square with a side of 60 mm is considered. The aggregate content is 0.5, $h = 0.03$ mm, and the Fuller gradation is adopted with sizes from 0.3 to 9.5 mm. The random walk paths are illustrated in Figure 9. Figure 9 shows that the Brownian particle first wanders in the lower left domain and then shifts to the upper left domain by crossing the bottom edge. After that, it explores the central domain for a while and finally comes near the upper right vertex.
Before applying the RWA to concrete, a reasonable value of $R_0$ needs to be known. Although a larger $R_0$ results in a more accurate $D_{\text{con}}$, a higher computational cost is required. In addition, the finite element size also leads to a slight fluctuation of $D_{\text{con}}$ for different initial locations of the Brownian particle. The two shortcomings could be overcome by applying the ergodic hypothesis \cite{35}, i.e., $D_{\text{con}}$ is taken as the average value over the random walks of $M$ Brownian particles. For this purpose, a cubic element with a side of 20 mm is selected. The Fuller gradation is adopted with sizes from 0.3 to 2 mm.

For such a situation, periodic boundary conditions are imposed. When it walks across the other faces, edges, or vertexes, the procedure can be performed in a similar manner. To demonstrate the detailed walk process of a Brownian particle, a two-dimensional simulation square with a side of 60 mm is considered. The aggregate content is 0.5, $h = 0.03$ mm, and $\mu = 2.5$. The maximum distance of the initial Brownian particle from the center of the simulation is taken as $R_0$. Based on the developed algorithm, a computer program is written with FORTRAN language for reconstructing the mesostructure of concrete and for implementing the random walk of Brownian particle. The two shortcomings could be overcome by applying the ergodic hypothesis \cite{35}, i.e., $D_{\text{con}}$ is taken as the average value over the random walk $s$ of $M$ Brownian particles. For this purpose, a cubic element with a side of 20 mm is selected. The Fuller gradation is adopted with sizes from 0.3 to 2 mm.
from 0.3 to 9.5 mm, \( f_a \) is 0.1, 0.3, 0.5, and 0.7, and \( D_i/D_{bcp} = 5 \). Thus, \( D_i/D_{bcp} \) is plotted against \( R_0 \) in Figure 10, which indicates that \( D_i/D_{bcp} \) first decreases gradually with the increase of \( R_0 \) and then keeps unchanged when \( R_0 \) exceeds 20 mm. By taking \( R_0 \) as 40 mm, \( D_i/D_{bcp} \) is plotted against \( M \) in Figure 11, which indicates that \( D_i/D_{bcp} \) fluctuates slightly for a smaller value of \( M \). When \( M \) exceeds 250, \( D_i/D_{bcp} \) keeps unchanged. The results are similar for other \( f_a \), \( D_i/D_{bcp} \), and aggregate gradations. As a conservative estimate, \( R_0 = 40 \) mm and \( M = 450 \) are adopted in this paper.

![Figure 10](image1.png)

**Figure 10.** Effect of \( R_0 \) on chloride diffusivity of concrete.

![Figure 11](image2.png)

**Figure 11.** Effect of \( M \) on chloride diffusivity of concrete.
As stated in the previous section, the EAM is constructed in this paper to reduce computational cost. To show the validity of the EAM, a computer simulation is performed on a cubic element with a side length of 20 mm on a Thermaltake workstation. In the simulation, the Fuller gradation is adopted with sizes from 0.3 to 9.5 mm, $f_a$ varies from 0.1 to 0.6, and $D_i/D_{bcp}$ is taken as 5, 10, and 15. The estimated $D_i/D_{bcp}$ is plotted in Figure 12, indicating that $D_i/D_{bcp}$ with EAM is very close to that without EAM. For a given value of $D_i/D_{bcp}$ at 5, 10, and 15, the average relative error between them is 0.858%, 1.88%, and 2.56%, respectively. However, the computational time is greatly reduced. For example, when the EAM is adopted, the computational time for concrete with $f_a = 0.5$ decreases by 72.6%. Therefore, the EAM gives similar results but at significantly reduced computational cost.

![Figure 12. Comparison of estimated chloride diffusivity of concrete with and without equivalent aggregate model (EAM).](image)

5. Experimental Verification

To verify the validity of the developed RWA, a chloride diffusion test was performed. In the test, specimens with a w/c of 0.6 were cast with ordinary Portland cement. The Fuller gradation was adopted with sizes from 0.3 to 9.5 mm and $f_a$ was 0.0, 0.15, 0.35, 0.55, and 0.75. After 24 h of casting, these specimens were demolded and cured in water at 21 °C for 28 days. The electrical conductivity method was adopted to measure $D_{con}$ [36]. The test device is schematically shown in Figure 13.
Since h was not measured, it has to be estimated. It was confirmed that, for normal-strength concrete, h is in the range between 0.01 to 0.05 mm [27]. Thus, h is set to be 0.03 mm, i.e., the average between 0.01 and 0.05 mm. In this test, $D_{bcp}$ (i.e., for concrete with zero aggregate volume fraction) was measured as $15.4 \times 10^{-12} \text{ m}^2/\text{s}$. At present, it is very tough to directly determine $D_i$ in the laboratory, and therefore this paper resorts to experimental calibration. In this test, $D_{con}$ at $f_a = 0.75$ was measured as $5.38 \times 10^{-12} \text{ m}^2/\text{s}$. Thus, $D_i$ is obtained, by the inverse method, as $62.5 \times 10^{-12} \text{ m}^2/\text{s}$. With these parameters known, the random walk is performed to estimate $D_{con}$, as shown in Figure 14, which shows a good agreement between the RWA and the measured $D_{con}$. The relative error between them is 4.99%, 0.70%, and 3.25% for $f_a = 0.15, 0.35,$ and 0.55, respectively.

To further verify the RWA, the experimental data of Zheng and Zhou [37] are chosen. In their experiment, the cement type, aggregate gradation, accelerated method, and curing conditions are the same as those in the last experiment. But w/c was 0.5 and $f_a$ was 0.0, 0.3, 0.4, 0.5, 0.6, and 0.7. The experientially measured $D_{con}$ was shown in Figure 15. Likewise, $h = 0.03 \text{ mm}$. In the test, $D_{con}$ was measured as $8.21 \times 10^{-12} \text{ m}^2/\text{s}$ and $4.29 \times 10^{-12} \text{ m}^2/\text{s}$ for $f_a = 0.0$ and 0.7, respectively. The former is $D_{bcp}$, i.e., $D_{bcp} = 8.21 \times 10^{-12} \text{ m}^2/\text{s}$ and the latter is adopted to calibrate $D_i$ which is obtained as $51.6 \times 10^{-12} \text{ m}^2/\text{s}$. Thus, $D_{con}$ is computed by the RWA, as shown in Figure 15, indicating that the RWA agrees well with the experimental results. When $f_a = 0.3, 0.4, 0.5,$ and 0.6, the relative error is 5.94%, 4.26%, 0.36%, and 4.08%, respectively.
Further, the experimental data of Yang and Su [8] are considered. In this experiment, mortar specimens with water/cement ratio 0.4 were cast with ASTM Type I Portland cement and $fa$ was 0.0, 0.1, 0.2, 0.3, and 0.4. The aggregate volume percentage $P_{v,j}$ passing the sieve with radius $R_j$ at 0.075, 0.15, 0.30, 0.59, 1.18, 2.375, and 4.75 mm is 2.83%, 11.4%, 40.3%, 68.1%, 91.6%, 99.7%, and 100%, respectively. After 24 h of casting, the mortar specimens were demolded and cured in water at 23 °C for 12 months. A migration method was used to determine $D_{con}$. As in the above two verification examples, $h = 0.03$ mm, $D_{bcp}$ was measured as $2.03 \times 10^{-12}$ m$^2$/s, and $D_i$ was calibrated as $10.2 \times 10^{-12}$ m$^2$/s from the value of $D_{con}$ measured at $fa = 0.4$. Thus, a comparison between the RWA and the experimental results is made as shown in Figure 16, which shows a good agreement between them. When $fa = 0.1$, 0.2, and 0.3, the relative error between them is 1.53%, 0.520%, and 1.90%, respectively. Therefore, the validation of the RWA is preliminarily verified.

Recently, Shafikhani and Chidiac [38] derived an approximate analytical solution of the chloride diffusivity of concrete based on phenomenological multi-scale models. By considering the effects of aggregate and ITZ on the diffusion of chloride ions in concrete separately, the solution can be expressed as

$$D_{con} = \frac{3(1 - fa)^2}{3 - fa} \cdot \frac{12 + 18hS_a}{12 - 9hS_a} D_{bcp}$$

(27)

where $S_a$ is the aggregate surface area per unit volume of concrete and given by

$$S_a = 4\pi(R^2)N_v f_a$$

(28)
With the size distribution of aggregates and $D_{bcp}$ known, the approximate analytical solution can be calculated as shown in Figure 16. It is seen from Figure 16 that, compared with the numerical method, the approximate analytical solution underestimates the experimental results and the deviation from the experimental results increases with the increase of $f_a$. The relative error between them is 0.941%, 3.09%, 10.7%, and 17.2% for a given $f_a$ at 0.1, 0.2, 0.3, and 0.4, respectively. The reason for this could be that the interactions between aggregates and ITZs are not fully considered. Therefore, the numerical method developed in this paper is of higher accuracy than the approximate analytical solution.

As seen in the previous sections, there are two limitations in this study. First, to increase the computational efficiency, the EAM is adopted. As a result, the connectivity of ITZs cannot be fully embodied and the accelerated effect is neglected. Second, two of data points from each experimental database have to be chosen to calibrate the chloride diffusivities of bulk cement paste and ITZ, thereby reducing the number of data points that are used for verification. Therefore, further experimental investigations with more data points need to be conducted to provide enough statistical evidence for the validity of the RWA.

When the chloride diffusivity of concrete is determined, Fick’s second law can be used to compute the chloride profiles in concrete analytically or numerically [39–41] and the time for the reinforcement surface to reach the critical chloride content [42,43], which is beyond the scope of this paper.

6. Conclusions

Based on the RWA, a numerical method has been proposed for evaluating the chloride diffusivity of concrete. The concrete mesostructure with periodic aggregate distribution has been reconstructed for a given aggregate volume fraction and gradation and an estimated ITZ thickness. When the chloride diffusivities of two sets of concrete samples are measured, the inverse method has been used to calibrate that of ITZ. To decrease the computing cost, an EAM has been constructed. The ITZ thickness and chloride diffusivity of the equivalent aggregate have been derived. It has been shown
that $h_{eq}/h$ decreases with increasing $f_a$ and $h$ and that, compared with the Fuller gradation, the EVF one has a smaller value of $h_{eq}/h$. $D_{eq}/D_i$ is almost directly and inversely proportional to $h_{eq}$ and $R$ respectively. With the EAM, concrete is simplified as a two-phase material. The chloride diffusivity has been estimated with the three-dimensional RWA with periodic boundary conditions on the movement of Brownian particles. The main advantage of the RWA is that there is no need to simulate the detailed zigzag walk of a Brownian particle with finite step sizes and therefore the computational time is greatly reduced and the sensitivity of chloride diffusivity to the step size is eliminated. Finally, the validation of the RWA has preliminarily been verified with experimental results.

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