A shape theorem for the scaling limit of the IPDSAW at criticality

Philippe Carmona
and Nicolas Pétrélis

Laboratoire de Mathématiques Jean Leray UMR 6629
Université de Nantes, 2 Rue de la Houssinière
BP 92208, F-44322 Nantes Cedex 03, France
e-mail: philippe.carmona@univ-nantes.fr
e-mail: nicolas.petrelis@univ-nantes.fr

Abstract: In this paper we give a complete characterization of the scaling limit of the critical Interacting Partially Directed Self-Avoiding Walk (IPDSAW) introduced in Zwanzig and Lauritzen (1968). As the system size $L \in \mathbb{N}$ diverges, we prove that the set of occupied sites, rescaled horizontally by $L^{2/3}$ and vertically by $L^{1/3}$ converges in law for the Hausdorff distance towards a non trivial random set. This limiting set is built with a Brownian motion $B$ conditioned to come back at the origin at the time at which its geometric area reaches 1. The modulus of $B$ up to a $\alpha_1$ gives the height of the limiting set, while its center of mass process is an independent Brownian motion.

Obtaining the shape theorem requires to derive a functional central limit theorem for the excursion of a random walk with Laplace symmetric increments conditioned on sweeping a prescribed geometric area. This result is proven in a companion paper Carmona and Pétrélis (2017).

MSC 2010 subject classifications: Primary 60K35; Secondary 82B26, 82B41.
Keywords and phrases: Polymer collapse, phase transition, shape theorem, Brownian motion, Local limit theorem.

Contents

1 Introduction and results ................................................................. 2
  1.1 The model ................................................................. 2
  1.2 Main result: the limiting shape of IPDSAW at criticality ................. 3
  1.3 Reminder: scaling limits in the non-critical regimes ...................... 4
2 Organization of the proof ............................................................ 5
  2.1 Random walk representation of IPDSAW at its critical point .......... 6
  2.2 Center-of-mass Walk and Profile ........................................ 6
  2.3 Outline of the proof of Theorem D ....................................... 9
3 Preparations ............................................................................. 12
  3.1 More about the renewal process ............................................ 12
  3.2 Skorohod’s representation Theorem for cadlag random functions .... 14
  3.3 Proof of Theorem B ..................................................... 15
4 Proof of Theorem A subject to Theorem D .................................. 15
  4.1 Tightness of $(y_L)_{L \in \mathbb{N}}$ (recall Remark 2.1) ...................... 15
  4.2 From the profile and the center-of mass walk to the occupied set, i.e., proof of Theorem A subject to Theorem C ................................................ 16
  4.3 Time change, i.e., proof of Theorem C subject to Theorem D ........ 16
    Outline of the proof of (4.6) ................................................. 17
5 Proof of Theorem D ................................................................. 20

†Acknowledgements. The author thanks the Centre Henri Lebesgue ANR-11-LABX-0020-01 for creating an attractive mathematical environment.

*July 5, 2018
1. Introduction and results

Deriving the scaling limit of a polymer model at its critical point is a difficult issue that had been tackled so far in Deuschel, Giacomin and Zambotti (2005) or in Sohier (2013) for wetting models and in Caravenna and Deuschel (2009) for a Laplacian pinning-model. With the present paper, we display the scaling limit of the critical-IPDSAW. It is a Shape Theorem whose limiting object is a truly 2-dimensional random set.

1.1. The model

The interacting partially directed self-avoiding walk (IPDSAW) is a self-avoiding random walk on $\mathbb{Z}^2$ that only takes unitary steps upwards, downwards and to the right. Thus, the set of allowed $L$-step paths is

$$W_L = \{w = (w_i)_{i=0}^L \in (\mathbb{N}_0 \times \mathbb{Z})^{L+1} : w_0 = 0, w_L - w_{L-1} = \rightarrow,$$

$$w_{i+1} - w_i \in \{\uparrow, \downarrow, \rightarrow\} \forall 0 \leq i < L - 1,$$

$$w_i \neq w_j \forall i < j\}.$$ Any non-consecutive vertices of the walk though adjacent on the lattice are called self-touchings and an energetic reward $\beta \geq 0$ is assigned to each trajectory for each self-touching. Thus, every random walk trajectory $w = (w_i)_{i=0}^L \in W_L$ is associated with the Hamiltonian

$$H_L(w) := \sum_{\substack{i,j=0 \\ i<j \rightarrow}}^L 1_{\{|w_i - w_j| = 1\}}, \tag{1.1}$$

which allows us to define $P_{L,\beta}$ the polymer law in size $L$ as,

$$P_{L,\beta}(w) = \frac{e^{\beta H_L(w)}}{Z_{L,\beta}}, \quad w \in W_L, \tag{1.2}$$

where $Z_{L,\beta}$ is the normalizing constant known as the partition function of the system. The exponential growth rate of the partition function is captured by the free energy of the model, i.e., $f(\beta) = \lim_{L \to \infty} \frac{1}{L} \log Z_{L,\beta}$.

The IPDSAW undergoes a collapse transition at some $\beta_c$ that is explicitly known (see e.g. Brak, Guttmann and Whittington (1992) or (Nguyen and Pétrélis, 2013, Theorem 1.3)) and the phase diagram is partitioned into an extended phase $\mathcal{E} = [0, \beta_c)$ inside which the free energy is larger than $\beta$ and a collapsed phase $\mathcal{C} = [\beta_c, \infty)$ where the free energy equals $\beta$. The asymptotics of the free energy close to criticality are analyzed in (Carmona, Nguyen and Pétrélis, 2016, Theorem B) where the phase transition is proven to be second order with a critical exponent $3/2$, i.e., $f(\beta_c - \epsilon) = \beta_c - \epsilon + \gamma \epsilon^{3/2} + o(\epsilon^{3/2})$ where the pre-factor $\gamma$ is closely related with a continuous model built with Brownian trajectories that are penalized energetically depending on their geometric area.

In Carmona, Nguyen and Pétrélis (2016) and Carmona and Pétrélis (2016), a rather complete description of the main geometric features of a typical path sampled from $P_{L,\beta}$ is provided inside the extended phase ($\beta < \beta_c$) and inside the collapsed phase ($\beta > \beta_c$) (see the discussion in Section 1.3 below). However, the scaling limit of the model at criticality ($\beta = \beta_c$), which is the most delicate case, was still to be derived and this is the object of the present paper.
1.2. **Main result: the limiting shape of IPDSAW at criticality**

We identify each \( w \in \Omega \) with a connected compact subset of \( \mathbb{R}^2 \) denoted by \( S(\omega) \) that extends the sites of \( \mathbb{Z}^2 \) occupied by \( w \) to squares of length 1, i.e.,

\[
S(w) = \left\{ \cup_{i=0}^L w(i) + \left[ -\frac{1}{2}, \frac{1}{2} \right]^2 \right\}, \quad w \in \Omega_L. \tag{1.3}
\]

For \((v_1, v_2) \in (0, \infty)^2\), we let \( T_{v_1,v_2} \) be a rescaling operator that acts on \( \mathcal{S} \) the set of closed subsets of \( \mathbb{R}^2 \) endowed with the Hausdorff distance. For \( S \subset \mathbb{R}^2 \), the set \( T_{v_1,v_2}(S) \) is obtained after rescaling \( S \) by \( v_1 \) horizontally and by \( v_2 \) vertically, i.e.,

\[
T_{v_1,v_2}(S) = \left\{ \left( \frac{x}{v_1}, \frac{y}{v_2} \right) : (x, y) \in S \right\}. \tag{1.4}
\]

For \((\alpha_1, \alpha_2) \in [0, 1] \), we denote by \( Q_{L,\beta}^{\alpha_1,\alpha_2} \) the law of \( T_{L,\alpha_1,L,\alpha_2}(S(w)) \) seen as a random variable on \( \mathcal{S} \) endowed with its Borel \( \sigma \)-algebra and when \( w \) is sampled from \( P_L, \beta \).

With Theorem A below, we prove that, at criticality, the IPDSAW rescaled in time by \( L^{2/3} \) and in space by \( L^{1/3} \) converges in distribution towards a non trivial random set, built with the help of two independent Brownian motions.

**Theorem A (Shape Theorem).** For \( \beta = \beta_c \), we have

\[
Q_{\beta}^{\frac{2}{3}, \frac{1}{3}} \xrightarrow{d} \mathcal{S}_{\text{crit}}(B, D) \tag{1.5}
\]

with \( \mathcal{S}_{\text{crit}}(B, D) \) a random subset of \( \mathbb{R}^2 \) defined as

\[
\mathcal{S}_{\text{crit}}(B, D) = \left\{ (x, y) \in [0, a_1] \times \mathbb{R} : D_x - \frac{|B_x|}{2} \leq y \leq D_x + \frac{|B_x|}{2} \right\} \tag{1.6}
\]

where \( B \) and \( D \) are independent Brownian motions of variance \( \sigma^2_\beta \) (defined below (2.1)), where \( a_1 \) is the time at which the geometric area described by \( B \) reaches 1, that is, \( \int_0^{a_1} |B_u| \, du = 1 \) and with \( B \) conditioned on the event \( B_{a_1} = 0 \).

Let us say a few words about the 4 main challenges that we faced to prove Theorem A. Thanks to the representation Theorem B, everything boils down to studying a random walk \( V \) conditioned on having a prescribed large geometric area. To be more specific, we need to consider the joint convergence of a couple of processes: the profile \( |V| \) (corresponding to \( |B| \) in Theorem A) and the center-of-mass walk \( M \) (corresponding to \( D \) in Theorem A).

1. Proving the convergence of time-changed discrete processes with an implicit time-change to corresponding time-changed continuous processes.
2. Handling the fluctuations of the center-of-mass walk \( M \) on the excursions of the profile \( |V| \). The main difficulty is that these are not independent at fixed time horizon, although we shall prove that they are asymptotically independent.
3. Extending the pioneering work of Denisov, Kolb and Wachtel (2015) to obtain local limit theorems for a 3 component process, i.e., an excursion of the profile conditioned on having a large extension, the associated center-of-mass walk and the geometric areas. This issue is settled in Carmona and Pétrélis (2017).
4. Adapting to our needs the reconstruction procedure introduced in Deuschel, Giacomin and Zambrotti (2005).
1.3. Reminder: scaling limits in the non-critical regimes

In the present section, we will explain why the shape Theorem stated above completes the picture of the scaling limit of IPDSAW initiated in Carmona, Nguyen and Pétrélis (2016) and Carmona and Pétrélis (2016). To that aim we need first to recall the stretch representation of the model, and then to associate with every configuration its profile and its center-of-mass walk, from which the occupied set in (1.3) can be reconstructed. With these tools in hand we will briefly recall the scaling limits obtained in Carmona, Nguyen and Pétrélis (2016) and Carmona and Pétrélis (2016) concerning the extended and the collapsed regime of IPDSAW. We will terminate this section by explaining why the critical regime (that is the object of the present paper) is more delicate than the others.

**Stretch description of a path.** There is a natural representation of any path in $\mathcal{W}_l$ as a collection of oriented vertical stretches separated by one horizontal step. Thus, we set $\Omega_l := \bigcup_{N=1}^{L} \mathcal{L}_{N,l}$, where $\mathcal{L}_{N,l}$ is the set of all possible configurations consisting of $N$ vertical stretches that have a total length $L$, that is

$$\mathcal{L}_{N,l} = \left\{ l \in \mathbb{Z}^N : \sum_{n=1}^{N} |l_n| + N = L \right\}.$$  

(1.7)

A one to one correspondence between $\Omega_l$ and $\mathcal{W}_l$ is obtained by associating with a given $l \in \Omega_l$, the path $w_l$ of $\mathcal{W}_l$ that starts at 0, takes $|l_1|$ vertical steps north if $l_1 > 0$ and south if $l_1 < 0$, then takes one horizontal step, then takes $|l_2|$ vertical steps north if $l_2 > 0$ and south if $l_2 < 0$ then takes one horizontal step and so on...

For $N \in \{1, \ldots, L\}$ and $l \in \mathcal{L}_{N,l}$, the Hamiltonian associated with $w_l$ can be rewritten as

$$H_l(w_l) = H_l(l_1, \ldots, l_N) = \sum_{n=1}^{N-1} (l_n \& l_{n+1})$$  

(1.8)

where

$$x \& y = \begin{cases} |x| \land |y| & \text{if } xy < 0, \\ 0 & \text{otherwise.} \end{cases}$$  

(1.9)

Thus, the polymer measure in (1.2) becomes

$$P_{L,\beta}(l) = \frac{e^{\beta H_l(w_l)}}{Z_{L,\beta}}, \quad l \in \Omega_l.$$  

(1.10)

We recall 1.3 and we denote by $S(l)$ the occupied set associated with any $l \in \Omega_l$ (i.e. $S(l) = S(w_l)$). We observe that $S(l)$ can be fully reconstructed with two auxiliary processes, i.e., the center-of-mass walk $M_l$ and the profile $|l|$. To be more specific, we associate with each $l \in \mathcal{L}_{N,l}$ the profile $|l| = (|l_i|)_{i=0}^{N+1}$ (with $l_{N+1} = 0$ by convention) and the center-of-mass walk $M_l = (M_{l,i})_{i=0}^{N+1}$ that links the middles of each stretch consecutively, i.e., $M_{l,0} = 0$ and

$$M_{l,i} = l_1 + \cdots + l_{i-1} + \frac{l_i}{2}, \quad i \in \{1, \ldots, N\},$$  

(1.11)
and $M_{l,N+1} = l_1 + \cdots + l_N$.

**Particularities of the critical regime.** A consequence of the fact that the occupied set $S(l)$ associated with $l \in \Omega_\beta$ can be recovered from its profile and center of mass walk is that for every $(\alpha_1, \alpha_2) \in [0, 1]^2$ the scaling limit of the rescaled occupied set $T_{l^{-1}, l^{-\infty}}(S(l))$ (with $l$ a typical path sampled from $P_{l, \beta}$) can be derived from the scaling limit of $(|l|, M_l)$ rescaled in time by $L^{\alpha_1}$ and in space by $L^{\alpha_2}$. This is the strategy adopted in Carmona, Nguyen and Pétrélis (2016) for the collapsed regime and in Carmona and Pétrélis (2016) for the extended regime. In the extended regime (i.e., $\beta < \beta_\star$), the horizontal extension of a typical path follows a law of large number of speed $L$ (that is $\alpha_1 = 1$), the vertical fluctuations of its center-of-mass walk are of order $\sqrt{L}$ (i.e., $\alpha_2 = 1/2$) whereas its vertical stretches are of finite size. Therefore, once rescaled vertically by $\sqrt{L}$ the profile vanishes whereas the center-of-mass walk displays a Brownian limit. In other words, once rescaled vertically by $\sqrt{L}$ and horizontally by $L$ and in the limit $L \to \infty$ the upper and lower envelopes of $T_{l, \sqrt{L}}(S(l))$ coalesce into a continuous trajectory whose law is that of a Brownian motion, i.e., one can straightforwardly deduce from (Carmona and Pétrélis, 2016, Theorem 2.8) that

$$Q_{L, \beta}^{1/2} \xrightarrow{d \quad L \to \infty} \{(s, \sigma_\beta B_s) : s \in [0, 1]\}$$

where $\sigma_\beta$ and $\epsilon_\beta$ are explicit constants and $B$ is a standard Brownian motion.

In the collapsed regime (i.e., $\beta > \beta_\star$) the fraction of self-touching performed by a typical trajectory equals $1 + o(1)$, which forces the vertical stretches to be long and with alternating signs. As a consequence the typical horizontal extension of a path sampled from $P_{L, \beta}$ is much shorter than its counterpart in the extended regime and follows a law of large number of speed $\sqrt{L}$ (i.e., $\alpha_1 = 1/2$). The typical length of vertical stretches is $\sqrt{L}$ as well (i.e., $\alpha_2 = 1/2$) and the profile rescaled in time and space by $\sqrt{L}$ converges towards a deterministic Wulff shape. The center-of-mass walk, in turn, fluctuates with an amplitude $L^{1/4}$ and therefore vanishes when we rescale it in time and space by $\sqrt{L}$. Unlike the extended regime, inside the collapsed phase the scaling limit of $T_{L, \sqrt{L}}(S(l))$ is driven by the profile only and we recall (Carmona, Nguyen and Pétrélis, 2016, Theorem D) which states that

$$Q_{L, \beta}^{1/2} \xrightarrow{d \quad L \to \infty} \mathcal{S}_{\beta}$$

where $\mathcal{S}_{\beta}$ is a deterministic Wulff-Shape, symmetric with respect to the $x$-axis.

In (Carmona and Pétrélis, 2016, Theorem 2.2) we proved that, at criticality, the horizontal extension of a typical path follows a central limit theorem with speed $L^{2/3}$ and a limiting law corresponding to that of the random time $\alpha_1$ at which the geometric area swept by a Brownian motion (of variance $\sigma_\beta^2$) reaches 1 conditioned on the fact that the Brownian touches 0 at $\alpha_1$. Thus, the last pending issue concerning the scaling limits of IPDSAW was to derive the scaling limit of the full path at criticality. This is the object of the present paper but let us insist on the fact that this is also the hardest issue. The reason is that, unlike the extended regime or the inside of the collapsed regime, at criticality the profile and the center-of-mass walk display vertical fluctuations of the same order (i.e. $L^{1/3}$).

2. **Organization of the proof**

The present Section is an extended outline of the proof of Theorem A. In Section 2.1 we settle some notation to state Theorem B which sheds light on the fact that the critical-IPDSAW can be studied indirectly with the help of an auxiliary random walk conditioned on sweeping a prescribed geometric area. Then, in Section 2.2 we state Theorem C which provides the scaling limits of the properly rescaled profile and center-of-mass walk for a typical configuration sampled from $P_{L, \beta}$. Theorem C actually implies Theorem A but we will not prove Theorem C directly. As exposed carefully in Remark 2.2, we will rather apply a time change on both profile and center-of-mass walk to state Theorem D which implies Theorem C but turns out to be easier to prove.
2.1. Random walk representation of IPDSAW at its critical point

The stretch representation of IPDSAW (displayed in Section 1.3 above) was initially used in Nguyen and Pétrélis (2013) to develop an alternative probabilistic approach of the model. This new approach involves an auxiliary random walk that we describe below before stating Theorem B which enlightens the particular relationship between this random walk and the model at criticality (i.e., at \( \beta = \beta_c \)).

We let \( P_\beta \) be the law of the random walk \( V := (V_n)_{n \in \mathbb{N}} \) starting from the origin and whose increments \((U_i)_{i \in \mathbb{N}}\) are i.i.d and follow a discrete Laplace law, i.e.,

\[
P_\beta(U_1 = k) = \frac{e^{-\frac{|k|}{\sigma}}}{\sigma} \quad \forall k \in \mathbb{Z} \quad \text{with} \quad \sigma_\beta := \frac{1+e^{a/2}}{1-e^{-a/2}}.
\]

and we set \( \sigma_\beta^2 := \text{Var}_\beta(U_1) \). For \( L \in \mathbb{N} \) and \( N \in \{1, \ldots, L\} \) we set

\[
\gamma_{N,L,N} := \{V: G_N(V) = L - N, V_{N+1} = 0\} \quad \text{with} \quad G_N(V) = \sum_{i=0}^{N} |V_i|,
\]

and we denote by \( T_N \) the one-to-one correspondence that maps \( \gamma_{N,L,N} \) onto \( \mathcal{L}_{N,L} \) as

\[
T_N(V)_i = (-1)^{i-1}V_i \quad \text{for all} \quad i \in \{1, \ldots, N\}.
\]

For \( n \in \mathbb{N} \) and for \( V = (V_i)_{i=0}^{\infty} \in \mathbb{Z}^\mathbb{N} \) we define \( K_n(V) := n + G_n(V) = \sum_{i=1}^{n} 1 + |V_i| \) and its pseudo-inverse

\[
\xi_s = \inf\{i \geq 0: K_i(V) \geq s\}, \quad s \in [0, \infty).
\]

We note incidentally that (2.4) implies \( K_j = \max\{n \geq 1: \xi_n = j\} \) for \( j \in \mathbb{N}_0 \). With a slight abuse of notation (and for random walk trajectories \( V \) only) we will call \( K_s \) the geometric area swept by \( V \) up to time \( n \) although it would be more correct to call it geometric area plus extension.

With these notations in hand, we state the fundamental Theorem B below. With this Theorem, we claim that at criticality, studying the model IPDSAW is completely equivalent to studying the \( V \) random walk conditioned on sweeping a prescribed geometric area.

**Theorem B** (Random Walk Representation at criticality).

\[
P_{L,\beta}(l \in \cdot) = P_\beta(T_{\xi_l}(V) \in \cdot | V_{\xi_l+1} = 0, K_{\xi_l} = L).
\]

This theorem will be proven in Section 3.3.

2.2. Center-of-mass Walk and Profile

With every \( l \in \Omega_L \), we associate \( \bar{M}_{l,L} \) and \( \bar{l}[l] \) the cadlag processes on \([0, \infty)\) obtained by rescaling the center-of-mass walk \( M_l \) and the profile \(|l|\) by \( L^{2/3} \) horizontally and by \( L^{1/3} \) vertically, i.e.,

\[
\bar{M}_{l,L}(s) = \frac{1}{L^{1/3}} M_{l,L}(L^{2/3} s \downarrow N_l), \quad s \in [0, \infty),
\]

\[
\bar{l}[l](s) = \frac{1}{L^{1/3}} |l| (L^{2/3} s \downarrow N_l), \quad s \in [0, \infty),
\]

where \( N_l \) is the number of vertical stretches composing \( l \) (i.e. \( l \in \mathcal{L}_{N_l,l} \)).

We denote by \( R_{L,\beta} \) the law of \((|l|_{L}, \bar{M}_{l,L})\) with \( l \) sampled from \( P_{L,\beta} \) and we state Theorem C which claims that the rescaled profile and center-of-mass walk of a typical configuration of the critical-IPDSAW converge simultaneously towards Brownian motions stopped at some particular random time. This Theorem is illustrated with Figure 1, where an exact simulation of the critical IPDSAW is provided in length \( L = 60000 \).
Theorem C. At criticality (β = βc) we have
\[ R_{L,β} \xrightarrow{d} \left( \{B_{s∧\alpha_1}, D_{s∧\alpha_1}\}_{s \in [0,\infty]} \right) \quad (2.8) \]
where B and D are independent Brownian motions of variance \( \sigma^2 \), where \( \alpha_1 \) is the time at which the geometric area described by B reaches 1, that is, \( \int_0^{\alpha_1} |B_u|du = 1 \) and with B conditioned on the event \( B_{\alpha_1} = 0 \).

We will prove in Section 4.2 that Theorem C implies Theorem A. For this reason, the target of the present paper will become to prove Theorem C, but let us first recall Theorem B which allows us to view \( R_{L,β} \) as the law of two other cadlag processes built with the V random walk. In this spirit, for a given random walk trajectory V, we define \( M = (M_i)_{i=0}^{\infty} \), the counterpart of the center-of-mass walk introduced in (1.11) as
\[ M_i = \sum_{j=1}^{i} (-1)^{j+1} \frac{U_j}{2}, \quad i \in \mathbb{N}, \quad (2.9) \]
and with \( M_0 = 0 \). We let \( \hat{V}_L \) and \( \hat{M}_L \) be the cadlag processes obtained after rescaling V and M by \( L^{2/3} \) in time and by \( L^{1/3} \) in space and stopped at \( \xi_L \) (recall (2.4)), that is,
\[ \hat{V}_L : [0, \infty) \to \mathbb{R} \quad \text{and} \quad \hat{M}_L : [0, \infty) \to \mathbb{R} \quad (2.10) \]
\[ s \to L^{-2/3} V_{(L^{2/3})\xi_L}, \quad s \to L^{-2/3} M_{(L^{2/3})\xi_L} \]
A consequence of Theorem B is that,
\[ R_{L,β} \xrightarrow{\text{Law}} (\hat{V}_L, \hat{M}_L) \quad (2.11) \]
with V sampled from \( P_\beta (\cdot \mid V_{\xi_L+1} = 0, K_{\xi_L} = L) \). In the proof of Theorem C (see Section 4.3 below), we will use the representation of \( R_{L,β} \) in (2.11).

Renewal structure We introduce a renewal structure which roughly consists of the excursions of the V random walk away from the origin and turns out to be a fundamental tool of our analysis. To that purpose, we define a sequence of stopping times \( (\tau_k)_{k \in \mathbb{N}} \) similar to ladder times by the prescription \( \tau_0 = 0 \) and
\[ \tau_{k+1} = \inf \{ i > \tau_k : V_{i-1} \neq 0 \text{ and } V_{i-1}V_i \leq 0 \}, \quad (2.12) \]
so that the length of the k-th excursions is given by
\[ \eta_k = \tau_k - \tau_{k-1} \quad (k \geq 1), \quad (2.13) \]
and the area swept
\[ \mathcal{A}_k = |V_{\tau_{k-1}}| + \cdots + |V_{\tau_k-1}| \quad (k \geq 1). \quad (2.14) \]
For each excursion we consider the sum of its length and its geometric area. For this reason we define the quantity \( X_i = \eta_i + \mathcal{A}_i \), for \( i \in \mathbb{N} \) and, with a slight abuse of notation, we will call \( X_i \) the geometric area swept by the i-th excursion. We set \( S_0 = 0 \) and \( S_n = X_1 + \cdots + X_n \) for \( n \geq 1 \) so that we can define \( X \) a random set of points on \( \mathbb{N}_0 \) as
\[ X = \{ S_n, n \in \mathbb{N}_0 \}. \quad (2.15) \]
We will also need to consider \( v_L \) the number of excursions that have been completed by V when its geometric area reaches \( L \), i.e., \( v_L := \max \{ i \geq 0 : S_i \leq L \} \).

Remark 2.1. It turns out that it is sufficient to prove Theorem C with V sampled from \( P_\beta (\cdot \mid L \in \mathbb{X}) \) instead of \( P_\beta (\cdot \mid V_{\xi_L+1} = 0, K_{\xi_L} = L) \). Understanding this last point requires to define, for every \( L \in \mathbb{N} \), the random variable \( Y_L \) which, for V sampled from \( P_\beta (\cdot \mid V_{\xi_L+1} = 0, K_{\xi_L} = L) \), records the length along which V sticks to the origin before \( \xi_L + 1 \), i.e.,
\[ Y_L := \max \{ k \geq 0 : V_{\xi_L-k+1} = \cdots = V_{\xi_L+1} = 0 \}. \]
We will prove in Section 4.1 that \((y_L)_{L \in \mathbb{N}}\) forms a tight family of random variables and therefore it is sufficient to consider the events \(y_1 = k\) for finitely many \(k\). Moreover, for \(k \in \mathbb{N}\), the event \(y_1 = k\) yields that \(L - k + 1 \in \mathbb{X}\) and \(V_{L-1+k+1} = 0\), so that the trajectory \((V_i)_{i=0}^{L-1+k}\) has for law \(P_{\beta} (\cdot | L - k + 1 \in \mathbb{X}, V_{L-1+k+1} = 0)\). We conclude by noticing that the symmetric Laplace distribution of the increments of \(V\) yields that \((V_i)_{i=0}^{L-1+k}\) and \((M_i)_{i=0}^{L-1+k}\) have the same law when \(V\) is sampled from \(P_{\beta} (\cdot | L - k + 1 \in \mathbb{X})\) as when \(V\) is sampled from \(P_{\beta} (\cdot | L - k + 1 \in \mathbb{X})\).

**Remark 2.2.** Our strategy to prove Theorem C is reminiscent of the strategy used in Deuschel, Giacomin and Zambotti (2005) to derive the scaling limit of a particular polymer model, i.e., the critical wetting model. To be more specific, the authors prove that, at criticality, a 1+1-dimensional L-step random walk pinned at an horizontal hard-wall, constrained to start and end at the wall and rescaled in time by \(L\) and in space by \(\sqrt{L}\) converges in distribution towards the modulus of a Brownian bridge. To achieve this result they display a smart reconstruction of the path under the polymer measure using the following 4 features of their model:

1. the Hamiltonian only depends on \(s_L\) where \(s_L := \{0, y_1, \ldots, y_{n_L} = L\}\) is the set of pinned sites of a given L-step random walk path. Moreover, under the critical polymer measure, \(\frac{1}{L} s_L\) converges in law towards \(s_{\infty} = \{s \in [0,1]: \beta_s = 0\}\) with \((\beta_s)_{s \in [0,1]}\) a Brownian bridge,
2. under the a priori random walk law \(P\) and once conditioned on \(s_L := \{0, y_1, \ldots, y_{n_L} = L\}\) the excursions of the random walk are independent and their respective length are prescribed by the inter-arrivals of \(s_L\),
3. a random walk excursion of length \(N\) rescaled in time by \(N\) and in space by \(\sqrt{N}\) converges in distribution towards a standard Brownian excursion,
4. under the critical polymer measure, the inter-arrivals of \(s_L\), i.e., \((y_{i+1} - y_i)_{i \geq 0}\) are almost surely finite and heavy-tailed random variables.

Their technique consists of using (i) in combination with Skohorod’s representation theorem to first sample \(s_L\) under the polymer measure and \(s_{\infty}\) such that \(\frac{1}{L} s_L\) converges almost surely towards \(s_{\infty}\). Then, with (iv) they claim that it suffices to consider finitely many inter-arrivals (and therefore excursions) of \(s_L\) to reconstruct a fraction of the path arbitrary close to 1. Finally, they use (ii-iii) in combination with Skohorod’s representation theorem to sample random walk excursions on the longest inter-arrivals of \(s_L\) and recover the modulus of a Brownian bridge.

Let us briefly describe the 4 features of IPDSAW (a–d) which can be substituted to (i–iv) above in order to adapt the reconstruction technic in our context:

a) Theorem B implies that the Hamiltonian of IPDSAW is somehow absorbed in the auxiliary random walk law \(P_{\beta} (\cdot | L \in \mathbb{X})\). The sequence of cumulated geometric areas \(X_L := \mathbb{X} \cap [0,L] = \{0,x_1, \ldots, x_{n_L}\}\) (recall (2.15)) plays the role of \(s_L\) and \(\frac{1}{L} X_L\) converges in distribution towards \(X_{\infty} := \{s \in [0,1]: B_s = 0\}\) where \(B\) is defined as in the statement of Theorem C,

b) under the random walk law \(P_{\beta}\) and once conditioned on \(X_L\), Proposition 3.1 below guarantees that the excursions (in modulus) are independent and their respective geometric area are prescribed by the inter-arrivals of \(X_L\),

c) with Theorem F, we claim that, once conditioned on sweeping a geometric area \(N\), a random walk excursion and its associated center-of-mass walk rescaled in time by \(N^{2/3}\) and in space by \(N^{1/3}\) converge towards a Brownian excursion normalized by its area and an independent Brownian motion,

d) Under \(P_{\beta}\), the inter-arrivals of \(X_L\), i.e., \((z_{i+1} - z_i)_{i \geq 0}\) are almost surely finite and heavy-tailed random variables.

Although the statements (a–d) constitute the skeleton of our path reconstruction (borrowed from Deuschel, Giacomin and Zambotti (2005)), adapting this method to the context of critical-IPDSAW raises two major additional challenges that are addressed in the present paper. The first difficulty comes from the fact that we consider simultaneously the random walk \(V_L\) and its associated center-of-mass walk \(M_L\). If the random
walk comes back very close to the origin at the end of its excursions this is not the case of the center-of-mass walk. Therefore, one needs informations on the position of \( M \). This requires to control the fluctuations of \( M \) on the "short" excursions of \( V \) and it will be the object of Propositions 2.4 and 2.5 in Section 2.3.

The second difficulty comes from the fact that the system size \( L \) provides a conditioning on the geometric area rather than on the length of the paths taken into account. To be more specific, in the wetting model the paths are constrained to complete their last excursion with a total length that equals \( L \) whereas in the present model there are no constraint on the total length but the paths must complete their last excursion with a total geometric area that equals \( L \). This is the reason why we perform below a time change on \( V \) instead of \( M \).

Operating this time change allows us to state Theorem D whose proof is simpler although it is equivalent to Theorem C.

We will need to perform the same type of time change for the standard Brownian motion \( B \), i.e., for \( s \geq 0 \) we denote by \( A \), the geometric area swept by \( B \) up to time \( s \), that is,

\[
A_s = \int_0^s |B_u| \, du.
\]

The continuity and strict monotonicity of \( A \) allows us to define \( a \) as the inverse of \( A \), i.e., \( A_s = s \) for \( s \in [0, \infty) \). For \( B \) and \( D \) are two independent Brownian motions, we define the continuous processes \( \tilde{B}, \tilde{D} \) as the Brownian counterparts of \( \tilde{V}, B \) and \( \tilde{M} \) as the Brownian counterparts of \( \tilde{V}, M \), respectively, i.e.,

\[
\tilde{B} : [0, \infty) \to \mathbb{R} \quad \text{and} \quad \tilde{D} : [0, \infty) \to \mathbb{R}
\]

\[
s \to B_{s/a_s} \quad \text{and} \quad s \to D_{s/a_s}
\]

We denote by \( \tilde{R}_{L,B} \) the law of \( (\tilde{V}, \tilde{M}) \) when \( V \) is sampled from \( \mathbb{P}_L(\cdot \mid L \in \mathbb{X}) \) and we state Theorem D which is the counterpart of Theorem C with \( \tilde{V}_L, \tilde{M}_L, \tilde{B}, \tilde{D} \) instead of \( \tilde{V}, \tilde{M}, \tilde{B}, \tilde{D} \). In Section 4.3, we will display an explicit link between those quantities with equations (4.4) and (4.5) and we will prove that Theorem C is a consequence of Theorem D.

**Theorem D.** For \( \beta > 0 \),

\[
\tilde{R}_{L,B} \xrightarrow{\ell \to \infty} \left( [B_{a_s}, D_{a_s}] \right)_{s \in [0,1]}
\]

where \( B \) and \( D \) are independent Brownian motions of variance \( \sigma^2_\beta \) and \( a \) is the inverse function of the geometric area swept by \( B \) and where \( B \) is considered under the conditioning \( \{B_{a_s} = 0\} \).

The proof of Theorem D is the object of Section 2.3 below.

### 2.3. Outline of the proof of Theorem D

Our proof of Theorem D relies on the renewal structure introduced in (2.12–2.15) above and it may be divided into three steps.
1. handling small excursions of random walk with Proposition 2.4, of Brownian motions with Proposition 2.5,
2. handling large excursions with Theorems E, F and G,
3. reconstructing the limiting process with Proposition 2.6.

Here, the geometric area of each excursion will be of particular importance. For \( k \in \mathbb{N} \), we will indeed truncate the rescaled profile \( |\tilde{V}_t| \) and the rescaled center-of-mass walk \( \tilde{M}_t \) (respectively the time-changed Brownian motions \( |\tilde{B}| \) and \( \tilde{D} \)) outside the excursions of \( V \) (resp. \( B \)) sweeping a geometric area larger than \( L/k \) (resp. \( 1/k \)) to obtain \( |\tilde{V}_{L,k}| \) and \( \tilde{M}_{L,k} \) (resp. \( |\tilde{B}| \) and \( \tilde{D} \)). Then, the proof of Theorem D will be organized as follows. With Proposition 2.4 (proven in Section 5.2), we state that provided \( k \) and \( L \) are large enough, \( ||\tilde{M}_{L} - \tilde{M}_{L,k}||_{\infty} + ||\tilde{V}_L - \tilde{V}_{L,k}||_{\infty} \) is arbitrary small in probability. With Proposition 2.5 (proven in Section 5.3) we prove that, provided \( k \) is large enough \( ||\tilde{B} - \tilde{B}_k||_{\infty} + ||\tilde{D} - \tilde{D}_k||_{\infty} \) also is arbitrarily small in probability. Finally, with Proposition 2.6 (proven in Section 5.1), we provide a simplified version of Theorem D by substituting the truncated processes \( (|\tilde{V}_{L,k}|, \tilde{M}_{L,k}) \) to \( (|\tilde{V}_L|, \tilde{M}_L) \) and \( (\tilde{B}_k, \tilde{D}_k) \) to \( (|\tilde{B}|, \tilde{D}) \), respectively. Those three propositions imply Theorem D.

**Remark 2.3.** For sake of conciseness, we will display the proof of Theorem D under the law \( P_{\beta, \mu_\beta} \) instead of \( P_\beta \). The only difference between those two laws is that, under \( P_{\beta, \mu_\beta} \) the law of \( V_0 \) is \( \mu_\beta \) (defined in (3.1)) which is symmetric on \( \mathbb{Z} \) with an exponential tail. Proving Theorem D under \( P_{\beta} \) is not more difficult but (and this is explained in Proposition 3.1 below) it would force us to consider separately the very first excursion of each path from all the other excursions. This distinction is not necessary anymore under \( P_{\beta, \mu_\beta} \) and this lightens the proofs a little bit.

**Truncation of the profile and center-of-mass walk.** We recall (2.9) and we observe that the center-of-mass walk can be written as

\[
M_t = -\frac{V_0}{2} + \sum_{j=1}^{i} (-1)^{j-1} \frac{V_j - V_{j-1}}{2} = \sum_{j=0}^{i-1} (-1)^{j-1} V_j + (-1)^{i-1} \frac{V_i}{2}, \quad i \in \mathbb{N}.
\]  

(2.21)

We recall (2.12) and for every \( r \in \mathbb{N} \), we let \( M^{\text{exc}}(r) \) be the contribution of the \( r \)-th excursion to the center-of-mass walk, i.e.,

\[
M^{\text{exc}}(r) = \sum_{i=\tau_{r-1}}^{\tau_r-1} (-1)^{i-1} V_i.
\]  

(2.22)

For \( x \in \mathbb{N} \), we truncate \( V \) outside the excursions of geometric area larger than \( x \) to obtain \( (V^+_x(i))_{i\in\mathbb{N}\cup\{0\}} \). Similarly, with the help of (2.22) we define the discrete process \( (M^+_x(i))_{i\in\mathbb{N}\cup\{0\}} \) which remains constant outside the excursions of geometric area larger than \( x \) and follows the center-of-mass walk elsewhere, i.e., for every \( t \in \mathbb{N} \) and \( i \in \{\tau_{r-1}, \ldots, \tau_r - 1\} \)

\[
M^+_x(i) := \sum_{r=1}^{\tau_r - 1} M^{\text{exc}}(r) 1_{\{X_i \geq x\}} + \left[ \sum_{j=\tau_{r-1}}^{\tau_r-1} (-1)^{j-1} V_j + (-1)^{r-1} \frac{V_r}{2} \right] 1_{\{X_i \geq x\}},
\]  

(2.23)

\[
V^+_x(i) := V_r 1_{\{X_i \geq x\}}.
\]

For \( k \in \mathbb{N} \), we let \( \tilde{V}_{L,k} \) and \( \tilde{M}_{L,k} \) be the cadlag processes obtained from \( V_{L,k} \) and \( M_{L,k} \) as we obtained \( \tilde{M}_L \) from \( M_L \), i.e.,

\[
\tilde{M}_{L,k}(s) := \frac{1}{L^{1/3}} M^+_{L,k}(s \xi_{L,k}), \quad s \in [0, 1],
\]  

(2.24)

\[
\tilde{V}_{L,k}(s) := \frac{1}{L^{1/3}} V^+_{L,k}(s \xi_{L,k}), \quad s \in [0, 1].
\]
**Proposition 2.4.** For every \( \epsilon > 0 \),

\[
\lim_{k \to \infty} \lim_{L \to \infty} \mathbb{P}_k \left( \sup_{u \in [0,1]} \left| \tilde{V}_L(s) - \tilde{V}_{L_k}(s) \right| \geq \epsilon \mid L \in \mathcal{X} \right) = 0,
\]

\[
\lim_{k \to \infty} \lim_{L \to \infty} \mathbb{P}_k \left( \sup_{u \in [0,1]} \left| \tilde{M}_L(s) - \tilde{M}_{L_k}(s) \right| \geq \epsilon \mid L \in \mathcal{X} \right) = 0,
\]

**Proposition 2.5.** For every \( \epsilon > 0 \),

\[
\lim_{k \to \infty} \mathbb{P}_k \left( \sup_{u \in [0,1]} \left| \tilde{B}(s) - \tilde{B}^k(s) \right| \geq \epsilon \right) = 0,
\]

\[
\lim_{k \to \infty} \mathbb{P}_k \left( \sup_{u \in [0,1]} \left| \tilde{D}(s) - \tilde{D}^k(s) \right| \geq \epsilon \right) = 0.
\]

where \( B \) and \( D \) are defined as in Theorem D.

Theorem E, proven for instance in Caravenna, Sun and Zygouras (Proposition A.8), claims that once rescaled by \( L \) and when sampled from \( \mathbb{P}_L(\cdot \mid L \in \mathcal{X}) \) the set \( \mathcal{X} \cap [0, L] \) converges in law (in the space of closed subsets of \([0, 1]\) endowed with the Hausdorff distance) towards \( \mathcal{C}_{1/3} := C_{1/3} \cap [0, 1] \) conditioned on \( 1 \in C_{1/3} \) where \( C_{1/3} \) is the 1/3-stable regenerative set.

**Theorem E.** For \( L \in \mathbb{N} \), we let \( \mathcal{X} \) be sampled from \( \mathbb{P}_L(\cdot \mid L \in \mathcal{X}) \), then,

\[
\lim_{L \to \infty} \frac{\mathcal{X} \cap [0, L]}{L} =_{\text{law}} \mathcal{C}_{1/3}.
\]

Theorems F and G below are proven in a companion paper Carmona and Pétrélis (2017). With Theorem F, we state that, a random walk excursion (together with its center-of-mass walk) conditioned to have a prescribed area \( L \), properly rescaled and subject to an adhoc time change converge towards a Brownian excursion normalized by its area (together with an independent Brownian motion) also subject to a similar time change.

**Theorem F.** We consider \( V \) sampled from \( \mathbb{P}_L(\cdot \mid X_1 = L) \), then,

\[
\lim_{L \to \infty} \left( \left| \tilde{V}_L(s), \tilde{M}_L(s) \right|_{s \in [0,1]} =_{\text{law}} \left( \mathcal{E}_a, B_a \right)_{s \in [0,1]} \right.
\]

where \( \mathcal{E} \) is a Brownian excursion normalized by its area and \( a \) is the inverse function of this area and where \( B \) is a standard brownian motion independent of \( \mathcal{E} \).
Let $Y$ be distributed as $(B_{n_1})_{0 \leq n_1 \leq 1}$ conditioned by $B_{n_1} = 0$. Then $Y$ is distributed as $\left(\frac{2}{\pi} \rho_1\right)^{2/3}$, $t \in [0, 1]$ where $(\rho_1)_{t \in [0,1]}$ is a Bessel bridge of dimension $\delta = 4/3$. We let $Z(Y)$ be the set of zeros of $Y$, i.e., $Z(Y) = \{s \in [0,1]; Y(s) = 0\}$. We let also $\gamma^Y$ be the law of an excursion of $Y$ renormalized by its extension and let $\gamma^d$ be the law of $\rho_\delta$ defined in (2.31) above.

**Theorem G.** The following equalities in distribution hold true,

$$Z(Y) = \text{Law } \tilde{C}_{1/3} \quad \text{and} \quad \gamma^Y = \gamma^d.$$  \hfill (2.32)

Propositions 2.4 and 2.5 are of key importance, because they reduce significantly the level of complexity of Theorem D. It becomes indeed sufficient to prove Proposition 2.6 below which is a simplified version of Theorem D to the extent that the profile and center of mass walk are replaced by their truncated version. We let $\tilde{R}_{L, \beta}^k$ be the joint law of $(\tilde{V}_{L,k}, \tilde{M}_{L,k})$ under $P_{\beta, \mu_\delta}$ $(\mid L \in \mathcal{X})$ and $B$ and $D$ be independent Brownian motions as defined in the statement of Theorem D.

**Proposition 2.6.** For every $k \in \mathbb{N}$,

$$\tilde{R}_{L, \beta}^k \xrightarrow{d} \left(\tilde{B}^k, \tilde{D}^k\right).$$  \hfill (2.33)

In section 5.1 we prove Proposition 2.6 subject to Theorems E, F and G.

3. Preparations

In Section 3.1 below, we give a complete description of the renewal structure introduced in (2.12–2.15) and consisting of excursions of the path away from the origin. We recall some facts from Carmona and Pétrélis (2016) concerning the geometric area and extension of those excursions and we go further by giving a method to reconstruct a trajectory $V$ of law $P_{\beta, \mu_\delta}$ with the help of independent excursions. In Section 3.2, we justify the use of Skorokhod Lemma for those cadlag processes considered in the present paper. In Section 3.3, we prove Theorem B.

3.1. More about the renewal process

We recall (2.12–2.15), we let $\mu_\beta$ be a probability law on $\mathbb{Z}$ defined as

$$\mu_\beta(k) = \frac{1 - e^{-\beta/2}}{2} e^{-\frac{k}{2\beta}} 1_{(k \neq 0)} + (1 - e^{-\beta/2})1_{(k = 0)},$$  \hfill (3.1)

and we let $P_{\beta, x}$ be the law of the random walk starting from $V_0 = x \in \mathbb{Z}$ and $P_{\beta, \mu_\delta}$ be the law of the random walk when $V_0$ has distribution $\mu_\beta$. In (Carmona and Pétrélis, 2016, Lemma 4.6) the tail distribution of $X_1$ is displayed as well as a renewal theorem for the set $\mathcal{X}$. To be more specific, for $x \in \{0, \mu_\beta\}$ there exists a $c_{x, \beta} > 0$ and $d_{x, \beta} > 0$ such that

$$P_{\beta, x}(X_1 = n) = c_{x, \beta} \frac{d_{x, \beta}}{n^{3/2}} (1 + o(1)) \quad \text{and} \quad P_{\beta, x}(n \in \mathcal{X}) = \frac{d_{x, \beta}}{n^{3/2}} (1 + o(1)).$$  \hfill (3.2)

Note that Lemma 4.6 is stated in Carmona and Pétrélis (2016) under $P_{\beta, \mu_\delta}$ but holds true under $P_{\beta}$ as well.

In the present paper we need to go further in the analysis of the renewal. With the help of $(\tau_k)_{k \geq 0}$, we divide any random walk trajectory $V$ into a sequence of excursions $(\xi_k)_{k \geq 0}$ and we also denote by $(|\xi_k|)_{k \geq 0}$ the same excursions in modulus, i.e., for $k \in \mathbb{N}$

$$\xi_k = (i, V_i)_{i \in (\tau_{k-1}, \tau_k)} \quad \text{and} \quad |\xi_k| = (i, |V_i|)_{i \in (\tau_{k-1}, \tau_k)}.$$  \hfill (3.3)

We will consider this sequence under $P_{\beta}$ and $P_{\beta, \mu_\delta}$. It is not true that the excursions themselves are independent because the sign of any excursion depends on the sign of the preceding excursion. However, when considered in modulus, those excursions are independent.
Proposition 3.1. Under $P_\theta$ the random processes $(|E_k|)_{k \geq 1}$ are independent and the sequence $(|E_k|)_{k \geq 2}$ is IID. The law of $|E_1|$ is that of the first excursion (in modulus) of a random walk of law $P_{\theta,0}$ and for $k \geq 2$ the law of $|E_k|$ is that of the first excursion (in modulus) of a random walk of law $P_{\theta,\mu_k}$.

Under $P_{\theta,\mu_k}$ the random processes $(|E_k|)_{k \geq 1}$ are IID. The law of $|E_1|$ is that of the first excursion (in modulus) of a random walk of law $P_{\theta,\mu_k}$.

Proof. We note that $V$ is a Markov chain, that $\tau_k$ is a stopping time (for $k \in \mathbb{N}$) and that for every $x \in \mathbb{Z}$ the law of $|E_1|$ under $P_{\theta,x}$ equals the law of $|E_1|$ under $P_{\theta,|x|}$. Therefore, the proof of Proposition 3.1 will be complete once we show (by induction) that for every $k \in \mathbb{N}$, the random variable $|V_{\tau_k}|$ is independent of the $\sigma$-algebra $\sigma(E_1, \ldots, E_k, \tau_k)$ and has the same law as $|T|$ with $T$ a random variable of law $\mu_\theta$.

We pick $t_1 < t_2 < \cdots < t_k \in \mathbb{N}$ and $(v_0, \ldots, v_{t_k-1}) \in \mathbb{Z}^{t_k}$ that are compatible with the event $M_k := \{E_1 = (v_0, v_{t_1-1}), \ldots, E_k = (v_{t_k-1}, v_{t_k-1})\}$.

We set $x \in \mathbb{N}_0$ and we compute $C := P_{\theta,\mu_\theta}(M_k \cap \{|V_{t_k}| = x\})$ as

$$C = P_{\theta,\mu_\theta}(V_0 = v_0, \ldots, V_{t_k-1} = v_{t_k-1}, V_{t_k} = -\text{sign}(v_{t_k-1})x) = P_{\theta,\mu_\theta}(V_0 = v_0, \ldots, V_{t_k-1} = v_{t_k-1})P_{\theta,\mu_\theta}(U_1 = x + |v_{t_k-1}|) = P_{\theta,\mu_\theta}(M_k)P_{\theta,\mu_\theta}(U_1 \geq |v_{t_k-1}|).$$

The ratio on the r.h.s. in (3.4) is equal to $(1 - e^{-\frac{\theta}{2}}) e^{-\frac{k}{2}} x$ which is exactly $P(|T| = x)$ when $T$ has law $\mu_\theta$. This completes the proof. $\square$

Random walk reconstruction. Proposition 3.1 will allow us to reconstruct a random walk of law $P_{\theta,\mu_\theta}$ with a sequence of independent excursions (in modulus). With Definitions 3.2 below we give the details of this construction.

Definition 3.2. Let $(e_i)_{i \in \mathbb{N}}$ be an i.i.d. sequence of symmetric Bernoulli trials taking values 1 and $-1$. Let also $(W_0, W_1, \ldots, W_{\tau_1})$ be the first excursion of a trajectory $W$ with law $P_{\theta,\mu_\theta}$. Independently from $(e_i)_{i \in \mathbb{N}}$, let $(V_i^j)_{i \in \{0, \ldots, \tau_1\}, j \geq 1}$ be a sequence of independent copies of $(W_0, W_1, \ldots, W_{\tau_1-1})$ and set $\tau_0 = 0$ and $\tau_j = \tau_{j-1} + \tau_1^j$ for $j \geq 1$. Finally, define $V$ as follows:

- $V_i = V_i^j$, $i \in \{0, \ldots, \tau_1 - 1\}$,
- for $j \geq 2$ if $V_0^j \neq 0$ then $V_i = -\text{sign}(V_{\tau_{i-1}}) |V_{\tau_{i-1}}|$ for every $i \in \{\tau_{j-1}, \ldots, \tau_j - 1\}$,
- for $j \geq 2$ if $V_0^j = 0$ then $V_i = e_j |V_{\tau_{i-1}}|$ for every $i \in \{\tau_{j-1}, \ldots, \tau_j - 1\}$.

The resulting stochastic process $V$ is a random walk of law $P_{\theta,\mu_\theta}$.

Remark 3.3. The construction in Definition 3.2 will be used in Section 5.2 and we note that, by construction, the sequence of signs $(e_i)_{i \in \mathbb{N}}$ is independent of the modulus of the trajectory $(|V_i|)_{i \in \mathbb{N}_0}$ and also independent of $(\tau_i)_{i \in \mathbb{N}_0}$.

With Definitions 3.4 below we display an alternative construction in order to generate for a random walk of law $P_{\theta,\mu_\theta}(\cdot \mid L \in \xi)$. This construction will be used in Section 5.1.

Definition 3.4. Let $(e_i)_{i \in \mathbb{N}}$ be an i.i.d. sequence of symmetric Bernoulli trials taking values 1 and $-1$. Recall (2.15) (and the definition of $v_1$ below) and independently from $(e_i)_{i \in \mathbb{N}}$ sample a random set

$$\mathcal{X} = \{0, X_1, \ldots, X_1 + X_2 + \cdots + X_{\tau_{i-1}} - L\}$$

under $P_{\theta}(\cdot \mid L \in \xi)$. Independently from $(e_i)_{i \in \mathbb{N}}$ and $\mathcal{X}$, sample for every $j \in \{1, \ldots, \tau_i\}$ the excursion $(V_i^j)_{i \in \{0, \ldots, \tau_i\}}$ with law $P_{\theta,\mu_\theta}(\cdot \mid X = X_j)$. Set $\tau_0 = 0$ and $\tau_j = \tau_{j-1} + \tau_1^j$ for $j \geq 1$. Finally, define $V$ as follows:
Thus, for \( V \in V^L_i \), \( i \in \{0, \ldots, \tau_1 - 1\} \),

- for \( j \geq 2 \) if \( V^L_0 \neq 0 \) then \( V_i = -\text{sign}(V_{\tau_{j-1}}) |V^L_{\tau_{j-1}}| \) for every \( i \in \{\tau_{j-1}, \ldots, \tau_j - 1\} \),
- for \( j \geq 2 \) if \( V^L_0 = 0 \) then \( V_i = \epsilon_j |V^L_{\tau_{j-1}}| \) for every \( i \in \{\tau_{j-1}, \ldots, \tau_j - 1\} \).

The resulting stochastic process \( V \) is a random walk of law \( P_{\beta,x} (\cdot | L \in \mathcal{X}) \).

3.2. Skorohod’s representation Theorem for cadlag random functions

Along the paper, we will often need to consider some piecewise constant cadlag processes defined either on \([0, \infty)\) or on \([0, 1]\). We will also need to consider the interpolated versions of such processes. To that aim we define two sets of functions, i.e., for \( I \in \{0, 1\} \), we let \((\mathcal{C}_I, d)\) be the set of continuous functions on \( I \), endowed with

\[
d(f, g) = \sum_{k=1}^{\infty} \frac{1}{2^k} \frac{||f - g||_{C^0[k, k+1]}}{1 + ||f - g||_{C^0[k, k+1]}},
\]

and similarly we let \((D_I, d)\) be the set of cadlag functions defined on \( I \) also endowed with the same distance. We recall that \((\mathcal{C}_I, d)\) is a Polish space whereas \((D_I, d)\) is not. Therefore, one can a priori not apply directly the Skorohod’s representation Theorem in \((D_I, d)\) (see (Billingsley, 2008, Theorem 6.7)).

Let us explain briefly below how this difficulty can be handled.

For \( n \in \mathbb{N} \), we will only consider functions \( F \) that are piecewise constant and cadlag, defined on \( I \in \{0, 1\} \), and such that all jumps of \( F \) occur at times belonging to \( I \cap \frac{1}{n} \). For such functions, we denote by \( F_{\text{int}} \) their interpolated version, that is,

\[ F_{\text{int}} : I \mapsto \mathbb{R} \]

\[
s \mapsto (1 - \{sn\}) F(\frac{\{sn\}}{n}) + \{sn\} F(\frac{\{sn+1\}}{n}). \quad (3.5) \]

We recall (2.1–2.4). All the cadlag processes considered in the rest of the paper are built with the increments \((U_i)_{i=1}^{\xi_L}\) (resp. \((U_i)_{i=1}^{\tau_1}\)) of a random walk \( V \) of law \( P_{\beta,x} (\cdot | L \in \mathcal{X}) \) (respectively \( P_{\beta,x} (\cdot | X_1 = L) \)) with \( x \in \{0, \mu_\beta\} \). Since \( \xi_L \leq L \) and \( \tau_1 \leq L \), it is useful to define for \( \alpha > 0 \)

\[ \mathcal{A}_{L,\alpha} := \{i \in \{1, \ldots, L\} : |U_i| \geq \alpha \log L\}. \quad (3.6) \]

With (2.1) and (3.2) we easily prove that there exists an \( \alpha > 0 \) such that for \( x \in \{0, \mu_\beta\} \)

\[
\lim_{L \to \infty} P_{\beta,x} [\mathcal{A}_{L,\alpha} | L \in \mathcal{X}] = 0 \quad \text{and} \quad \lim_{L \to \infty} P_{\beta,x} [\mathcal{A}_{L,\alpha} | X_1 = L] = 0. \quad (3.7)
\]

As a consequence, if we denote by \( F_L \) (respectively \( F^\text{int}_L \)) a generic cadlag random process (and its interpolated version) built with the increments of \( V \) and rescaled vertically by \( L^\alpha \) for some \( \alpha > 0 \), we can deduce from (3.7) that for \( x \in \{0, \mu_\beta\} \) and for \( \varepsilon > 0 \)

\[
\lim_{L \to \infty} P_{\beta,x} [d(F_L, F^\text{int}) > \varepsilon | L \in \mathcal{X}] = 0 \quad \text{and} \quad \lim_{L \to \infty} P_{\beta,x} [d(F_L, F^\text{int}) > \varepsilon | X_1 = L] = 0. \quad (3.8)
\]

Thus, for \( F^\infty \in \mathcal{C}_I \), the convergence in law of \((F_L)_{L \in \mathbb{N}} \) towards \( F^\infty \) in \((D_I, d)\) is equivalent to the convergence in law of \( F^\text{int}_L \) towards \( F^\infty \) in \((\mathcal{C}_I, d)\). Moreover, the fact that \( F_L \) only jumps at times belonging to \( I \cap \frac{1}{2} \) allows us to reconstruct \( F_L \) from \( F^\text{int}_L \) in an easy way. Therefore, Skorohod’s representation Theorem can be applied in the present paper for convergence in \((D_I, d)\) as well.
3.3. Proof of Theorem B

We recall the stretch description of IPDSAW in (1.7–1.9) and we observe that the partition function can be rewritten under the form

$$Z_{L,\beta} = \sum_{N=1}^{L} \sum_{l \in \mathcal{L}_{N,L}} e^{\beta \sum_{n=1}^{N} (l_{n+1} - l_{n})}. \tag{3.9}$$

We note that \(\forall x, y \in \mathbb{Z}\) one can write \(x - y = \frac{1}{2} (|x| + |y| - |x + y|)\) and therefore the partition function in (3.9) becomes

$$Z_{L,\beta} = \sum_{N=1}^{L} \sum_{l \in \mathcal{L}_{N,L}} \exp\left(\beta \sum_{n=1}^{N} |l_{n}| - \frac{\beta}{2} \sum_{n=0}^{N} |l_{n} + l_{n+1}|\right) = c_{\beta} e^{\beta L} \sum_{N=1}^{L} \prod_{l \in \mathcal{L}_{N,L}} \prod_{n=0}^{N} \exp\left(-\frac{\beta}{2} |l_{n} + l_{n+1}|\right). \tag{3.10}$$

At this stage we recall the definition of the auxiliary random walk \(V\) in (2.1–2.2) as well as the family of one to one correspondence \((T_N)_{N \geq 0}\) between path configurations and random walk trajectories (see 2.3). Since for \(l \in \mathcal{L}_{N,L}\) the increments \((U_{l})_{l=1}^{N+1}\) of \(V = (T_N)^{-1}(l)\) in (2.3) necessarily satisfy \(U_{l} := (-1)^{-1}(l_{1} - 1)\), one can rewrite (3.10) as

$$Z_{L,\beta} = c_{\beta} e^{\beta L} \sum_{N=1}^{L} \prod_{l \in \mathcal{L}_{N,L}} P_{\beta}(V_{N,L} = N) \quad \text{with} \quad \Gamma_{\beta} := \frac{c_{\beta}}{e^{\beta L}}. \tag{3.11}$$

The probabilistic representation of the partition function in (3.11) is a key tool when studying IPDSAW. It allows for instance to spot quickly the critical point of the model which turns out to be the solution in \(\beta = 1\). The present paper being fully dedicated to the critical regime of IPDSAW, we will henceforth always work at \(\beta = 1\) and therefore we remove the term \(\Gamma_{\beta}\) from the r.h.s. in (3.11).

Another useful consequence of formula (3.11) is that it provides us with a very strong link between the polymer law \(P_{L,\beta}\) and the random walk law \(P_{\beta}\) conditioned on a suitable event. We recall (2.4–2.4), the fact that \(\Gamma_{1} = 1\) and also that the term indexed by \(N\) in the sum in (3.10) corresponds to the contribution to the partition function of those path in \(\mathcal{L}_{N,L,N}\). Consequently, we can derive from (3.10–3.11) that for every \(N \in \{1, \ldots, L\}\),

$$P_{L,\beta}(N_{1} = N) = P_{\beta}(\xi_{1} = N | V_{\xi_{1}+1} = 0, K_{\xi_{1}} = L),$$

$$P_{L,\beta}(l \in \cdot | N_{1} = N) = P_{\beta}(T_{\beta}(V) \in \cdot | \xi_{1} = N, V_{N+1} = 0, K_{N} = L). \tag{3.12}$$

Theorem B is a straightforward consequence of (3.12).

4. Proof of Theorem A subject to Theorem D

4.1. Tightness of \((Y_{k})_{k \in \mathbb{N}}\) (recall Remark 2.1)

Let \(k_{0} \in \mathbb{N}\), recall that \(P_{\beta}(U_{1} = 0) = 1/c_{\beta}\) and observe that

$$P_{\beta}(Y_{L} > k_{0}, V_{\xi_{L}+1} = 0, K_{\xi_{L}} = L) = \sum_{k=k_{0}+1}^{L} \left(\frac{1}{c_{\beta}}\right)^{k} P_{\beta}(V_{\xi_{L}+1} = 0, V_{\xi_{L}+1} = 0, K_{\xi_{L}+1} = L-k+1)$$

$$= \sum_{k=k_{0}+1}^{L} \left(\frac{1}{c_{\beta}}\right)^{k} P_{\beta}(L-k+1 \in \mathcal{X}, V_{\xi_{L}+1} = 0)$$

$$= (1 - \frac{1}{c_{\beta}}) \sum_{k=k_{0}+1}^{L} \left(\frac{1}{c_{\beta}}\right)^{k} P_{\beta}(L-k+1 \in \mathcal{X}). \tag{4.1}$$
Note that (4.1) can also be written without the event \( \{ y_L > k_0 \} \) in the left hand side in (4.1) and with a sum running from \( k = 0 \) to \( k = L \) in the r.h.s. Therefore, we recall (3.2) and we obtain that there exists a \( d > 0 \) such that \( P_R( V_{L+1} = 0, K_{L+1} = L ) \leq \frac{d}{L^m} (1 + o(1)) \). As a consequence (since also \( c_\beta > 1 \)) we deduce from (4.1) that for every \( \epsilon > 0 \) we can choose \( k_0 \) large enough such that

\[
P_R( y_L > k_0 \mid V_{L+1} = 0, K_{L+1} = L ) \leq \epsilon \quad \text{for every } L \in \mathbb{N}.
\]

4.2. From the profile and the center-of mass walk to the occupied set, i.e., proof of Theorem A subject to Theorem C

**Proof.** By the Skorohod's representation Theorem, we can already assert that there exists \((h_L, m_L)_{L \in \mathbb{N}}\) a sequence of cadlag processes and \( B \) and \( D \) two independent Brownian motions of variance \( \sigma_\beta^2 \), all defined on the same probability space \((\Omega, \mathcal{F}, P)\) so that

- \( P \)-a.s. it holds that for all \( K > 0 \),
  \[
  \lim_{L \to \infty} \sup_{s \in [0, K]} \left| (h_L(s), m_L(s)) - \left( |B_{3L^\alpha_1}|, D_{3L^\alpha_1} \right) \right| = 0, \quad (4.2)
  \]

- for all \( L \in \mathbb{N} \), \((h_L, m_L)\) has for law \( R_{L, \beta} \).

Then, we recall the definition of \( \mathcal{S}_{\text{crit}}(B, D) \) in (1.6) and we note that the random set

\[
S_L(h_L, m_L) := \left\{ (s, y) \in [0, i_L] \times \mathbb{R} : m_L(s) - \frac{|h_L(s)|}{2} \leq y \leq m_L(s) + \frac{|h_L(s)|}{2} \right\},
\]

(4.3) has for law \( Q_{L, \beta}^{3/4} \) with \( i_L = \sup \{ s \geq 0 : \int_0^s |h_L(s)| ds \leq 1 \} \).

Then, it remains to note that \( S_L(m_L, h_L) \) converges \( P \)-a.s. towards \( \mathcal{S}_{\text{crit}}(B, D) \) for the Hausdorff distance. The almost sure convergence of \((m_L, h_L)\) towards \((B_{\lambda \alpha_1}, D_{\lambda \alpha_1})\) implies that \( i_L \) also converges towards \( a_1 \) almost surely and this is sufficient to conclude.

**Remark 4.1.** To be completely rigorous, we must note that, as it is defined in (4.3), the law of \( S_L(h_L, m_L) \) is not exactly \( Q_{L, \beta}^{3/4} \). However, we recall (1.3) and (1.4) and we observe that by enlarging \( S_L(h_L, m_L) \) of \( 1/L^{1/3} \) vertically and by shifting it of \( 1/2 L^{2/3} \) horizontally we retrieve a set of Law \( Q_{L, \beta}^{3/4} \). Finally, since the Hausdorff distance between those two sets is bounded above by \( 1/L^{1/3} \), working with \( S_L(h_L, m_L) \) is sufficient to conclude.

4.3. Time change, i.e., proof of Theorem C subject to Theorem D

As explained in Remark 2.1, \( \tilde{V}_L \) is considered under the conditioning \( \{ L \in X \} \) whereas \( \tilde{B} \) is considered under the conditioning \( \{ B_{\lambda \alpha_1} = 0 \} \). We observe that for \( s \geq 0 \),

\[
\tilde{V}_L(s) = \frac{K_{L, \alpha_1 \beta}}{L} \quad \text{and} \quad \tilde{B}_L = \tilde{B}(A_{\lambda \alpha_1}),
\]

and similarly

\[
\tilde{M}_L(s) = \frac{K_{L, \alpha_1 \beta}}{L} \quad \text{and} \quad \tilde{D}_L = \tilde{D}(A_{\lambda \alpha_1}).
\]

Therefore, by (Billingsley, 2008, Lemma p. 151), the proof of Theorem C will be complete once we show that the following convergence in law holds true

\[
\lim_{L \to \infty} \left( \frac{\tilde{V}_L}{\tilde{M}_L}, \frac{K_{L, \alpha_1 \beta}}{L} \right) =_{\text{law}} \left( |\tilde{B}|, \tilde{D}, A_{\lambda \alpha_1} \right).
\]

(4.6)
The following relations between $\xi'$ and $\bar{V}$ on the one hand and between $a$ and $\bar{B}$ on the other hand will be of key importance to get (4.6)

$$a_s = \int_0^s \frac{1}{|B_{a(u)}|} du = \int_0^s \frac{1}{|B(u)|} du \quad \text{and} \quad \xi_{sL} = \int_0^s \frac{L}{1 + |V_{\xi_s}|} du = \int_0^s \frac{L^{2/3}}{L^{-1/3} + |V_{\xi_s}(u)|} du$$

(4.7)

where the first equality holds true for $s \in [0, 1]$ and the second for $s$ in the set $J_L$ of hopping times of $s \to \xi_{sL}$.

Outline of the proof of (4.6)

We will follow the scheme below

1. With the help of Theorem D, we infer the Skorohod’s representation Theorem and state that there exists a sequence of cadlag processes $|\bar{V}_L|$, $\bar{M}_L'$ and $|\bar{B}|$, $\bar{D}'$ defined on the same probability space $(\Omega, \mathcal{F}, \mathbb{P})$ such that for $P$-a.e. $\omega \in \Omega$

$$\lim_{L \to \infty} \left\| |\bar{V}_L'| - |\bar{B}'| \right\|_{\infty, [0,1]} = 0 \quad \text{and} \quad \lim_{L \to \infty} \left\| \bar{M}_L' - \bar{D}' \right\|_{\infty, [0,1]} = 0$$

(4.8)

and such that $(|\bar{V}_L'|, \bar{M}_L')$ has the same law as $(|\bar{V}_L|, \bar{M}_L)$ under the conditioning $\{L \in \chi\}$ and $(|\bar{B}', \bar{D}')$ are two independent Brownian motions of variance $\sigma^2$ under the conditioning $\{B_{a_1} = 0\}$.

2. We define for $s \in [0, 1]$, the quantity $a'(s)$ with the l.h.s. of formula 4.7 applied to $\bar{B}'$ and for $L \in \mathbb{N}$ the quantity $\xi_{sL}'$ with the r.h.s. of (4.7) applied to $\bar{V}_L'$. Then we show

**Lemma 4.2.** For all $\varepsilon > 0$,

$$\lim_{L \to \infty} \mathbb{P} \left( \sup_{s \in [0,1]} \left| \frac{\xi_{sL}'}{L^{2/3}} - a'(s) \right| \geq \varepsilon \right) = 0.$$  

(4.9)

3. Subsequently, we define the quantity $A'$ as the inverse of $a'$ and $K'$ as

$$K'(j) = \max \{ i \geq 1 : \xi_s' = j \}, \quad j \leq \xi_{sL}'$$

and we show the following convergence in probability.

**Lemma 4.3.** For all $\varepsilon > 0$,

$$\lim_{L \to \infty} \mathbb{P} \left( \sup_{\xi \in [0,\infty]} \left| \frac{K_{\varepsilon,\left|L^{2/3}\right|\xi_s, \left|L^{2/3}\right|\xi_s}}{L} - A'(\xi_s, 1) \right| \geq \varepsilon \right) = 0.$$  

(4.10)

4. At this stage, (4.8) and Lemma 4.3 allow us to state that

$$\lim_{L \to \infty} \left( |\bar{V}_L'|, \bar{M}_L', \frac{K_{\varepsilon,\left|L^{2/3}\right|\xi_s, \left|L^{2/3}\right|\xi_s}}{L} \right) \Rightarrow \text{law} \left( |\bar{B}'|, \bar{D}', A'(1) \right)$$

(4.11)

and this implies (4.6) by a straightforward application of (Billingsley, 2008, Lemma p. 151).

**Remark 4.4.** We note that, since we defined $\xi'$ with the help of formula (4.7), it is a continuous process and therefore it does not have the same law as $\xi$ as defined in (2.4). However this difference is armless because the set of times $J_s'$ at which $\xi'$ takes integer values has the same law as the set $J_s$ containing the hopping times of $\xi$ and moreover between two consecutive points of $J_s$ (respectively $J_s'$) $\xi$ (resp. $\xi'$) jumps by one unit exactly.

At this stage it remains to prove lemmas 4.2 and 4.3 and we begin with the proof of (4.9).
Proof. We recall (4.7). The proof of (4.9) will be complete once we show that for every $\varepsilon > 0$,

$$
\lim_{L \to \infty} \mathbb{P}\left( \int_0^1 \frac{1}{L^{-1/3} + |V'_L(u)|} - \frac{1}{|B'_L(u)|} \, du \geq \varepsilon \right) = 0.
$$  

(4.12)

We define for $L \in \mathbb{N}$ and $\eta > 0$ the three quantities

$$
C_1 = \int_0^1 \frac{1}{L^{-1/3} + |V'_L(u)|} - \frac{1}{|B'_L(u)|} \, du, 
$$  

(4.13)

$$
C_2 = \int_0^1 \frac{1}{|B'_L(u)|} \, du, 
$$  

(4.14)

$$
C_3 = \int_0^1 \frac{1}{L^{-1/3} + |V'_L(u)|} \, du.
$$  

(4.15)

We immediately observe that (4.8) and the dominated convergence theorem yield that for $\eta > 0$ and for $\mathbb{P}$-a.e $\omega \lim_{L \to \infty} C_1 = 0$. Moreover, (4.7) combined with the fact that $B$ and $B'$ have the same law and with the fact that $q_1 < \infty$ yields that for $\mathbb{P}$-a.e $\omega$ the function $u \mapsto 1/|B'_u|$ is integrable on $[0,1]$. This is sufficient to conclude that for $\mathbb{P}$-a.e $\omega$, $\lim_{\eta \to 0} C_2 = 0$.

Thus it remains to consider $C_3$. Since $V'_L$ and $V_L$ are equally distributed, it comes that

$$
C_3 = \frac{1}{L} \int_0^L \frac{L^{1/3}}{1 + |V_{\xi(x)}|} \, dx \leq \frac{1}{L^{2/3}} \left| \{ j \leq \xi_L : |V_j| \leq \eta L^{1/3} \} \right|.
$$  

(4.16)

Therefore, the proof of (4.12) will be complete once we show that for all $\varepsilon > 0$

$$
\lim_{L \to \infty} \limsup_{\eta \to 0} \mathbb{P}_\beta \left( \left| \left\{ j \leq \xi_L : |V_j| \leq \eta L^{1/3} \right\} \right| \geq \varepsilon L^{2/3} \mid L \in \mathcal{X} \right) = 0.
$$  

(4.17)

To prove (4.17), we use some results obtained in Carmona and Pétrélis (2016) under the same conditioning. We recall (2.12–2.15) and we denote by $(X^{(i)})_{i=1}^{N_k}$ the order statistics of $(X_i)_{i=1}^{N_k}$ and by $(\eta L^{1/3})_{i=1}^{N_k}$ the sequence of horizontal excursions reordered according to the sequence $(X^{(i)})_{i=1}^{N_k}$. Then, we distinguish between the $k$ largest such excursions and their lengths, i.e., $(X^{(i)})_{i=1}^{k}$ and the others, i.e.,

$$
\frac{1}{L^{2/3}} \left| \left\{ j \leq \xi_L : |V_j| \leq \eta L^{1/3} \right\} \right| \leq A_{k,L}^\eta + B_{k,L}.
$$  

(4.18)

with

$$
A_{k,L}^\eta = \frac{1}{L^{2/3}} \sum_{j=1}^{k} \left| \left\{ i \in \left\{ \tau_{r_j-1}, \ldots, \tau_{r_j} - 1 \right\} : |V_i| \leq \eta L^{1/3} \right\} \right| \quad \text{and} \quad B_{k,L} = \sum_{i=k+1}^{N_k} \eta L^{2/3}. \quad (4.19)
$$

In the second step of the proof of (Carmona and Pétrélis, 2016, Proposition 4.7), it is shown that for all $\varepsilon > 0$

$$
\lim_{k \to \infty} \limsup_{L \to \infty} \mathbb{P}_\beta (B_{k,L} \geq \varepsilon L \mid L \in \mathcal{X}) = 0. \quad (4.20)
$$

Thus, the proof of (4.17) will be complete once we show that for all $k \in \mathbb{N}$

$$
\lim_{\eta \to 0} \limsup_{L \to \infty} \mathbb{P}_\beta \left( \left| \left\{ i \in \left\{ \tau_{r_j-1}, \ldots, \tau_{r_j} - 1 \right\} : |V_i| \leq \eta L^{1/3} \right\} \right| \geq \varepsilon L^{2/3} \mid L \in \mathcal{X} \right) = 0. \quad (4.21)
$$

which again will be a consequence of the fact that for all $j \in \mathbb{N}$

$$
\lim_{\eta \to 0} \limsup_{L \to \infty} \mathbb{P}_\beta \left( \left| \left\{ i \in \left\{ \tau_{r_j-1}, \ldots, \tau_{r_j} - 1 \right\} : |V_i| \leq \eta L^{1/3} \right\} \right| \geq \varepsilon L^{2/3} \mid L \in \mathcal{X} \right) = 0. \quad (4.22)
$$
From (Carmona and Pétrélis, 2017, Proposition 3.4) we know that $\tau/\ell$ and (4.22) will be proven once we show that

$$R_{\ell,\eta,\epsilon}(j) := \mathbb{P}_\beta\left(\{i \in \{\tau_{j-1}, \ldots, \tau_j - 1\} : |V_i| \leq \eta L^{1/3}\} \mid L \in \mathcal{X}\right) \geq \epsilon L^{2/3} \mid L \in \mathcal{X}$$

(4.23)

$$= \sum_{\ell=1}^{J_j} \mathbb{P}_\beta(X_{j,\ell} = \ell \mid L \in \mathcal{X}) \max_{x \in \{0, \mu_\epsilon\}} \mathbb{P}_\beta, x\left(\{i \in \{0, \ldots, \tau_1 - 1\} : |V_i| \leq \eta L^{1/3}\} \mid L \in \mathcal{X}, X_1 = \ell\right).$$

We can argue here that $X_j$ under $\mathbb{P}_\beta(\cdot \mid L \in \mathcal{X})$ converges in distribution towards the $j$-th largest inter-arrival of an 1/3-stable regenerative set on $[0, 1]$ conditioned on 1 being in the set. For this reason, and for every $\xi > 0$ there exists an $m_1 > 0$ such that for $L$ large enough

$$\mathbb{P}_\beta(X_{j,\ell} \notin [m_1, L] \mid L \in \mathcal{X}) \leq \xi,$$

(4.24)

Thus, we let

$$C_{L,\eta,\epsilon} := \{V : \{i \in \{0, \ldots, \tau_1 - 1\} : |V_i| \leq \eta L^{1/3}\} \geq \epsilon L^{2/3}\}$$

and (4.22) will be proven once we show that

$$\lim_{\eta \to 0} \limsup_{L \to \infty} \sup_{\ell \in [m_1, L]} \max_{x \in \{0, \mu_\epsilon\}} \mathbb{P}_\beta, x(C_{L,\eta,\epsilon} \mid X = \ell) = 0.$$  

(4.25)

From (Carmona and Pétrélis, 2017, Proposition 3.4) we know that $\tau/\ell^{2/3}$ under $\mathbb{P}_{\beta,\mu_\epsilon}(\cdot \mid X = \ell)$ (or under $\mathbb{P}_\beta(\cdot \mid X = \ell)$) converges in distribution towards $\mathcal{R}$ the extinction of a Brownian excursion normalized by its area. Thus, for every $\xi > 0$, there exists $[\alpha_1, \alpha_2] \subset (0, \infty)$ such that for $L$ large enough, we state that for every $\ell \in [m_1, L]$ we have

$$\max_{x \in \{0, \mu_\epsilon\}} \mathbb{P}_\beta, x\left(\frac{\tau}{\ell^{2/3}} \in [\alpha_1, \alpha_2] \mid X = \ell\right) \leq \xi.$$  

As a consequence, (4.23) will be proven once we show that

$$\lim_{\eta \to 0} \limsup_{L \to \infty} \sup_{\ell \in [m_1, L]} \max_{x \in \{0, \mu_\epsilon\}} \mathbb{P}_\beta, x(C_{L,\eta,\epsilon} \cap \left\{\frac{\tau}{\ell^{2/3}} \in [\alpha_1, \alpha_2]\right\} \mid X = \ell).$$

(4.26)

At this stage, we introduce the notation $(\tilde{V}_j)_{j \in [0, \ell]} = (\frac{1}{\sqrt{m_1}} |V_j|)_{j \in [0, \ell]}$ and we note that for $\ell \in [m_1, L]$ and $\tau \in [\alpha_1, L^{2/3}, \alpha_2 L^{2/3}]$ we have

$$C_{L,\eta,\epsilon} \cap \left\{\frac{\tau}{\ell^{2/3}} \in [\alpha_1, \alpha_2]\right\} \subset \left\{\int_0^1 1\left\{\frac{1}{\sqrt{m_1}} |\tilde{V}_s| \leq \frac{\eta}{\sqrt{m_1}}\right\} ds \geq \frac{\epsilon}{\alpha_2}\right\},$$

(4.27)

where we have used that $L^{1/3} \leq \frac{\sqrt{\tau}}{\sqrt{m_1}}$ and $L^{2/3} \geq \frac{\tau}{\alpha_2}$. At this stage, we use the convergence established in (Carmona and Pétrélis, 2017, Theorem A) and (4.27) to assert that

$$\limsup_{L \to \infty} \sup_{\ell \in [m_1, L]} \max_{x \in \{0, \mu_\epsilon\}} \mathbb{P}_\beta, x(C_{L,\eta,\epsilon} \cap \left\{\frac{\tau}{\ell^{2/3}} \in [\alpha_1, \alpha_2]\right\} \mid X = \ell) \leq \mathbb{P}\left(\int_0^1 1\left\{\frac{1}{\sqrt{m_1}} |\tilde{V}_s| \leq \frac{\eta}{\sqrt{m_1}}\right\} ds \geq \frac{\epsilon}{\alpha_2}\right)$$

(4.28)

and we conclude by observing that $\mathbb{P}$ almost surely $\mathcal{E}$ is continuous on $[0, \mathcal{R}_\epsilon]$ and equals 0 at 0 and $\mathcal{R}_\epsilon$ only. Thus, the r.h.s. in (4.28) vanishes as $\eta \to 0$ and this concludes the proof.

At this stage, it remains to prove (4.10).
Proof. Let us define
\[ D_{L, \varepsilon} := \left\{ \sup_{t \in (0, \infty)} \frac{\sum_{i=0}^{L} A_i - A_{\varepsilon}}{L} \geq \varepsilon \right\} \]
and for \( \eta > 0 \) we set
\[ E_{L, \eta} := \left\{ \sup_{t \in [0, 1]} \left| \frac{\xi_{\varepsilon} \{ \xi_{\varepsilon} \geq \varepsilon \}}{L^{2/3}} - g(\varepsilon) \right| \leq \eta \right\}. \]

By lemma 4.2, the proof of (4.10) will be complete once we show that there exists \( \eta > 0 \) such that
\[ \lim_{L \to \infty} P(D_{L, \varepsilon} \cap E_{L, \eta}) = 0. \]
Under the event \( D_{L, \varepsilon} \cap E_{L, \eta} \) there exists a \( c \in [0, a_1 \wedge \xi_{\varepsilon}/L^{2/3}] \) such that
\[ \frac{1}{L} K_{L^{2/3}} - A_c \geq \varepsilon. \]
Assume that \( K_{L^{2/3}} \geq (A_c + \varepsilon)L \) (the other case is treated similarly). Then \( c L^{2/3} \geq \\
\xi_{L^{2/3}} \geq \xi_{(A_c + \varepsilon)L} \) and under \( E_{L, \eta} \) we can state that \( \xi_{(A_c + \varepsilon)L} \geq (a_{A_c + \varepsilon} - \varepsilon)L^{2/3} \) so that finally \( c \geq a_{A_c + \varepsilon} - \eta \).

But since \( a_c = c \) we get \( \eta \geq \min\{a_{u+} - a_u : u \in [0, 1 - \varepsilon]\} \). Since \( a \) is \( P \)-almost surely continuous and strictly increasing on \([0, 1]\) we complete the proof of the Lemma by claiming that
\[ \lim_{\eta \to 0} P(\min\{a_{u+} - a_u : u \in [0, 1 - \varepsilon]\} \leq \eta) = 0. \]

\[ \square \]

5. Proof of Theorem D

5.1. Truncated version of Theorem D, i.e., proof of Proposition 2.6

We recall Definition 3.4 and for every \( L \in \mathbb{N} \), we will generate a random walk path \((V_t)_{t=0}^{\xi_L} \) of law \( \mathbb{P}_L(\cdot | L \in \mathcal{X}) \). To begin with, we use Theorem E in combination with Skorohod’s representation theorem (the set of closed subsets in \([0, 1]\) endowed with the Hausdorff distance being a Polish space) to assert that there exists a sequence of random sets \( \mathcal{X}_L \) and a random set \( \mathcal{X}_\infty \) defined on the same probability space \((\Omega_1, \mathcal{F}_1, \mathbb{P}_1)\) and such that

1. For every \( L \in \mathbb{N} \), \( \mathcal{X}_L \) has the same law as \( \mathcal{X} \cap [0, L] \) with \( \mathcal{X} \) sampled from \( \mathbb{P}_L(\cdot | L \in \mathcal{X}) \). We will denote by \((X^L_j)_{j=1}^n\) the inter-arrivals of \( \mathcal{X}_L \), i.e.,
\[ \mathcal{X}_L = \{0, X^L_1, \ldots, X^L_j, \ldots, X^L_{n-1}, L\}, \]
2. \( \mathcal{X}_\infty \) is a \( C_{1/3} \) regenerative set intersected with \([0, 1]\) and conditioned on \( 1 \in \mathcal{X} \),
3. \( \lim_{L \to \infty} \frac{1}{L} \mathcal{X}_L(\omega) = \mathcal{X}_\infty(\omega) \) for \( \mathbb{P}_1 \)-a.e. \( \omega \).

For \( k \in \mathbb{N} \) and every \( \omega_j \in \Omega_1 \) we denote by \((d^\omega_j, f^\omega_j)\), \((d^\omega^\infty_j, f^\omega^\infty_j)\) the positions of the maximal intervals of \([0, 1]\) which are not intersecting \( \mathcal{X}_\infty \) and are larger than \( 1/k \). Note that \( r := r^\infty(\omega) \) is random and bounded above by \( k \). Note also that \( X^\omega_j := f^\omega_j - d^\omega_j > 1/k \) for every \( j \in \{1, \ldots, r\} \). We also denote by \( \tau_L \) the number of intervals of \([0, L]\) larger than \( L/k \) whose extremities are consecutive points of \( \mathcal{X}_L \). We let \((d^L_j, f^L_j)_{j=1}^{\tau_L}\) be the extremities of those intervals and \((X^L_{j+1})_{j=1}^{\tau_L}\) their associated inter-arrivals (i.e., \( 1 \leq e_1, e_2, \ldots, e_{\tau_L} \leq \tau_L \)). Because of the almost sure convergence of \( \frac{1}{L} \mathcal{X}_L \) towards \( \mathcal{X}_\infty \), we can claim that for \( L \) large enough \( \tau_L \) equals \( r \\) and moreover that for \( \mathbb{P}_1 \)-a.e. \( \omega_1 \) and \( j \in \{1, \ldots, r\} \)
\[ \lim_{L \to \infty} (d^L_j, f^L_j) = (d^\infty_j, f^\infty_j) \quad \text{and} \quad \lim_{L \to \infty} X^L_{e_j, e_j} = X^\infty_j. \] (5.1)

We will also need the notations
\[ \Lambda^\infty = \bigcup_{j=1}^{\tau_L} [d^\infty_j, f^\infty_j] \quad \text{and} \quad \Lambda_L = \bigcup_{j=1}^{\tau_L} [d^L_j, f^L_j]. \]

At this stage, we sample a family of independent random variables \((Y^{1,N}_{j,N})_{(j,N) \in \mathbb{N}} \) on a probability space \((\Omega_2, \mathcal{F}_2, \mathbb{P}_2)\) as follows
1. for every \((j, N) \in \mathbb{N}^2\) the random variable \(Y_{j, N}\) is a Bernoulli with parameter \(P_{\beta, \mu} (V_0 \neq 0 | X = N)\).

At this stage we use Theorem F and the Skorohod’s representation theorem to define on the same probability space \((\Omega_3, \sigma_3, P_3)\) a family of independent sequences of discrete random processes \(\{\tilde{e}^{i,j}_N\}_{N \in \mathbb{N}, \tilde{e} \in \{0, 1\}}\) and a family of independent continuous random processes \(\{e^{i,j}_N (s), D^{i,j}_N (s)\}_{s \in [0, a(1)], \tilde{e} \in \{0, 1\}, j \in \mathbb{N}, \tilde{e} \in \{0, 1\}}\) such that

1. for every \((j, N) \in \mathbb{N}^2\) the random process \(\tilde{e}^{i,j}_N = (V_i^{j,0,N})_{i=0}^{\tau_j,N}\) has for law \(P_{\beta, \mu} (\cdot | V_0 = 0, X = N)\), and the random process \(e^{i,j}_N = (V_i^{j,1,N})_{i=0}^{\tau_j,N}\) has for law \(P_{\beta, \mu} (\cdot | V_0 \neq 0, X = N)\).

2. for every \((j, y) \in \mathbb{N} \times \{0, 1\}\), \(\tilde{e}^{i,j}\) is a Brownian excursion normalized by its area and \(a_{j,y}\) is the inverse function of this area and \(D_{j,y}\) is a standard Brownian motion independent of \(\tilde{e}^{i,j}\).

3. for every \((j, y) \in \mathbb{N} \times \{0, 1\}\) and for \(P_3\)-a.e. \(\omega_3\)

\[
\lim_{N \to \infty} \sup_{s \in [0, 1]} \frac{1}{N^{1/3}} \left| V_{\tau_{j,N}}^{j,y,N} - \mathcal{E}^{i,j,y} (a_{j,y} (s)) \right| = 0 \quad \text{and} \quad \lim_{N \to \infty} \sup_{s \in [0, 1]} \frac{1}{N^{1/3}} \left| M_{\tau_{j,N}}^{j,y,N} - D_{j,y} (a_{j,y} (s)) \right| = 0.
\]

(5.2)

where \(\mathcal{E}^{i,j,y}_{N}\) is the pseudo-inverse of the geometric area (plus extension) of \(\mathcal{E}^{i,j,y}_N\) defined as in (2.4) and where \(M_{\tau_{j,N}}^{j,y,N}\) is the center-of-mass walk associated with \(\mathcal{E}^{i,j,y}_N\), i.e.,

\[
\begin{align*}
M_{\tau_{j,N}}^{j,y,N} &= \sum_{i=0}^{\tau_{j,N}-1} (-1)^{i+1} |V_i^{j,y,N}| + (-1)^{i-1} \frac{|V_i^{j,y,N}|}{2}, & i \in \{0, \ldots, \tau_{j,N} - 1\}, \\
M_{\tau_{j,N}}^{j,y,N} &= \sum_{i=0}^{\tau_{j,N}-1} (-1)^{i+1} |V_i^{j,y,N}|.
\end{align*}
\]

(5.3)

Let us note that for every \((j, N) \in \mathbb{N}^2\), the random process \(\mathcal{E}^{i,j,y}_{N}\) is an excursion of law \(P_{\beta, \mu} (\cdot | X = N)\). To lighten the notations we will drop the \(L\) dependency of \((X^L_t, \alpha^L, \beta^L, f^L)\) and \(e_{j,L}\) when there is no risk of confusion.

With these tools in hand, we apply for every \(L \in \mathbb{N}\) the construction of \(V\) (given in Definition 3.4) with law \(P_{\beta, \mu} (\cdot | L \in \mathcal{X})\) on a probability space \((\times_{l=1}^\infty \Omega_l, \sigma_l, \mu_l, P_3)\). Of course \(\mathcal{X}\) plays the role of \(\mathcal{X}\), then for every \(t \in \{1, \ldots, V_t\}\) we sample on \((\Omega_3, \sigma_3, P_3)\) an excursion \(V^L\) of law \(P_{\beta, \mu} (\cdot | X = X_L)\) and independently an i.i.d. sequence of Bernoulli trials \(e_{j,L}\). We do so except that for the indices \((e_{j,L})_{j=1}^\infty\) we replace each excursion \(V_{\mathcal{E}^{j,L} (\omega_3)} (\omega_3)\) by the excursion \(\mathcal{E}^{j,L} (\omega_3)\) defined above.

For every \(s \in [0, 1]\) we set \(j_{s} := \min \{ t \leq r_L : f_t \geq s L \}\) and for \(j \in \{1, \ldots, r_L\}\) we let \(\alpha_{j,L}\) be the product of \((-1)^{r_{j-1}}\) with the sign of the excursion of \(V\) indexed by \(e_{j}\), i.e.,

\[
\begin{align*}
\alpha_{j,L} &= (-1)^{r_{j-1}} \text{sign}(V_{r_{j-1}}) = (-1)^{r_{j-1}} \left( 1_{\{V_{r_{j-1}} = 0\}} e_{j} - 1_{\{V_{r_{j-1}} = 0\}} \right) \text{sign}(V_{r_{j-1}}) \\
&= (-1)^{r_{j-1}} \left( 1_{\{V_{r_{j-1}} = 0\}} e_{j} - 1_{\{V_{r_{j-1}} = 0\}} \right) \text{sign}(V_{r_{j-1}}).
\end{align*}
\]

(5.4)

For the ease of notations, for \(t \leq k\) and \(u \in [0, \infty)\), we will use the following shortcuts in the computations below:

\[
\tilde{V}^t := V^{e_{j,L} \mathcal{E}^{j,L} (\omega_3)}_{\mathcal{E}^{j,L} (\omega_3)} \quad \text{and} \quad \tilde{M}^t := M^{e_{j,L} \mathcal{E}^{j,L} (\omega_3)}_{\mathcal{E}^{j,L} (\omega_3)}.
\]

(5.5)

We recall (2.22–2.24) and (5.3) and we observe that the truncated processes \(\tilde{M}_{L,k}\) and \(\tilde{V}_{L,k}\) obtained
from \( V \) can be written as
\[
|\tilde{V}_{t,k}(s)| = 1_{\Lambda_{\varepsilon}}(sL) \frac{1}{L^{1/3}} \left| V_{t,k} \right| = 1_{\Lambda_{\varepsilon}}(sL) \frac{1}{L^{1/3}} \left| \tilde{V}_{t,k} \right|
\]
\[
= 1_{\Lambda_{\varepsilon}}(sL) \left( \left[ \frac{X_{s\varepsilon}}{L} \right]^{1/3} \left( \left[ \frac{1}{X_{s\varepsilon}} \right]^{1/3} \left| \tilde{V}_{t,k} \right| \right) \right), \tag{5.6}
\]
and
\[
\tilde{M}_{t,k}(s) = \frac{1}{L^{1/3}} \left[ \sum_{i=1}^{N} M_{i,k}^{\text{exc}}(\varepsilon_{i}) + 1_{\Lambda_{\varepsilon}}(sL) \left( \sum_{i=1}^{N} (-1)^{i-1} V_{i} + (-1)^{i-1} \varepsilon_{i} \right) \frac{1}{L} \right] \tag{5.7}
\]
\[
= \sum_{i=1}^{N} \left[ \frac{X_{s\varepsilon}}{L} \right]^{1/3} \alpha_{i,k} \left( \left[ \frac{1}{X_{s\varepsilon}} \right]^{1/3} \tilde{M}_{i,k}^{\varepsilon} \right) + 1_{\Lambda_{\varepsilon}}(sL) \left( \frac{X_{s\varepsilon}}{L} \right)^{1/3} \alpha_{i,k} \left( \left[ \frac{1}{X_{s\varepsilon}} \right]^{1/3} \tilde{M}_{i,k}^{\varepsilon} \right). \tag{5.7}
\]

To complete the proof of Proposition 2.6 it remains to identify the limiting distribution of \( (\tilde{V}_{t,k}, \tilde{M}_{t,k}) \) (defined in in (5.6–5.7)). The tension of \( (\tilde{V}_{t,k}, \tilde{M}_{t,k}) \) is ensured first by the fact that \( r_{k} \leq k \) for every \( L \in \mathbb{N} \) and then by the fact that for every \( (j,y) \in \{1,\ldots,k\} \times \{0,1\} \) the convergences in (5.2) ensure that for \( N \geq N_{j,y}(\omega_{2}) \) the modulus of continuity of \( \frac{1}{N} \sum_{i=1}^{N} V_{i}^{j,y} \), \( \frac{1}{N} \sum_{i=1}^{N} M_{i,k}^{j,y} \) are arbitrarily small.

We conclude by saying that with high probability \( L/k \) is larger than \( \max_{1 \leq j \leq k \leq N(0,1)} N_{j,y} \). Therefore, we need to obtain the limiting law of finite dimensional distributions of \( (\tilde{V}_{t,k}, \tilde{M}_{t,k}) \) as \( L \to \infty \). To that aim, we define two auxiliary processes \( \tilde{B}_{t,k}(s) \) and \( \tilde{D}_{t,k}(s) \) as
\[
[\tilde{B}_{t,k}(s) = 1_{\Lambda_{\varepsilon}}(s) \left( X_{s\varepsilon}^{1/3} \xi_{k}(a_{j,k}^{s}(\frac{\xi_{k}^{\text{exc}}}{X_{s\varepsilon}^{1/3}})), \; s \in [0,1], \tag{5.8}
\]
where for every \( t \leq r \), we set \( Y_{t} := Y_{t}(\omega_{1}, \omega_{2}) = Y_{t} \left( X_{s\varepsilon}^{1/3} \right) \) (which actually explains the \( L \)-dependency of \( [\tilde{B}_{t,k}(s)] \)) and
\[
[\tilde{D}_{t,k}(s) = \sum_{i=1}^{N} \left( X_{s\varepsilon}^{1/3} \xi_{k}(a_{i,k}^{s}(1)), \; s \in [0,1]. \tag{5.9}
\]

We will complete the proof by first observing that (5.1–5.2) implies that for every \( s \in [0,1] \) the following convergences occurs for \( \Phi_{i}^{j} = \Phi_{i}^{j} \), a.e. \((\omega_{1})_{i} \in \mathbb{R}^{d} \) and \( t_{i} \in \mathbb{R} \), i.e.,
\[
\lim_{L \to \infty} \left| \tilde{V}_{t,k}(s) \right| = 0 \quad \text{and} \quad \lim_{L \to \infty} \left| \tilde{M}_{t,k}(s) \right| = 0
\]
and then by showing that for every \( L \in \mathbb{N} \), the two dimensional process \( (\tilde{B}_{t,k}, \tilde{D}_{t,k}) \) has the same distribution as that of \( (\tilde{B}_{t,k}^{1}, \tilde{D}_{t,k}^{1}) \) defined in the statement of Proposition 2.6. To prove this last point we prove below that we do not change the law of \( (\tilde{B}_{t,k}, \tilde{D}_{t,k}) \) by removing the terms \( (a_{i,k}^{s}(1))_{i=1}^{N} \) in (5.9).

The resulting process \( (\tilde{B}_{t,k}, \tilde{D}_{t,k}) \) does not depend on \( L \) anymore and a straightforward consequence of Theorem G ensures that this process is distributed as \( (\tilde{B}^{1}, \tilde{D}^{1}) \).

For \( t \in \mathbb{N} \) we let \( \mathcal{F}_{t} \) be the sub-\( \sigma \)-algebra of \( \mathcal{A}_{\varepsilon} \) defined as
\[
\mathcal{F}_{t} = \sigma(\xi_{i}^{s}, \xi_{i}^{s}, \mathcal{D}^{s}, \mathcal{D}^{s}, j \leq t, \; y \in \{0,1\}, \; N \in \mathbb{N}), \tag{5.11}
\]

The proof will be complete once we show that for every \( t \leq r \) the law of the
\[
\left( \xi_{t}^{s}(a_{i,k}^{s}(\frac{\xi_{k}^{\text{exc}}}{X_{s\varepsilon}^{1/3}})), \; a_{i,k}^{s}(1) \right)_{i \in \{0,1\}^{1/3}} \tag{5.12}
\]
conditioned on the sub-$\sigma$-algebra $\mathcal{A}_1 \otimes \mathcal{A}_2 \otimes \mathcal{F}_{t-1} \otimes \mathcal{A}_4$ equals the law of
\[
(\mathcal{E}'(a\left(\frac{t-d_n}{X}\right)), D(a\left(\frac{t-d_n}{X}\right)))_{t \in [0,d_n^s+X^\infty_n]}
\] (5.13)
with $\mathcal{E}$ a Brownian excursion normalized by its area, $a$ the inverse function of this area and $D$ is a standard brownian motion independent of $\mathcal{E}$. Such an equality indeed allows us to compute the characteristic function of any finite dimensional distribution of $([\tilde{B}_k],D^k_{t-1})$ by conditioning successively on $\mathcal{F}_{t-1}$ from $t = r$ up to $t = 1$, getting rid at each step of the random variable $\alpha_{t,L}$.

To prove this later equality in law we note first that the law of $\mathcal{E}_{t,L} = \sum_{j \in \mathbb{N}} M_{t \vee \_}^{\max} \left(\frac{t-d_n}{X}\right)$ conditioned on $\mathcal{A}_1 \otimes \mathcal{A}_2 \otimes \mathcal{F}_{t-1} \otimes \mathcal{A}_4$ does not depend on $Y_j(\omega_1, \omega_2)$ and equals the law of (5.13), second that the random variable $\alpha_{t,L}$ is $\mathcal{A}_1 \otimes \mathcal{A}_2 \otimes \mathcal{F}_{t-1} \otimes \mathcal{A}_4$-measurable and takes values $-1$ and $1$ only, third that for any $c \in \{-1, 1\}$ the laws of $(\mathcal{E}, D)$ and $(\mathcal{E}', cD)$ are equal.

5.2. The center-of-mass walk outside large excursions, i.e., proof of Proposition 2.4

Proposition 2.4 contains two limits. We will display the proof of the second limit only since the proof of the first limit is way easier. To be more specific, the first limit gives some control on the fluctuations of the $V$ random walk sampled from $P_{\beta,\mu,k} (\cdot \mid L \in \mathcal{X})$ outside its largest excursions (in terms of geometric area swept). The second limit is much more involved, essentially because, despite the $V$ random walk, the center-of-mass walk does not come back close to the origin at the end of every excursion of $V$.

We recall (2.23) and we observe that for every $s \in [0, 1]$ we have
\[
\tilde{M}_t(s) - \tilde{M}_{L,k}(s) = \frac{1}{L^{1/3}} \left( M_{t,L} - M_{L,k}^{\max}(\xi_L) \right)
\] (5.14)

Therefore, Proposition 2.4 will be proven once we show that for every $\eta > 0$
\[
\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta,\mu,k} \left( \max_{n \leq \xi_L} \left| M_{t,L} - M_{L,k}^{\max}(\xi_L) \right| \geq \eta L^{1/3} \mid L \in \mathcal{X} \right) = 0.
\] (5.15)

Step 1: Proof of (5.15) subject to Claims 5.1 and 5.2

We recall (2.22) and for every $j \in \mathbb{N}$ we set
\[
T_{j,k} := \sum_{i=1}^{j} M_{t \vee \_}^{\max} (r) \mathbb{1}_{\{s, X \leq \xi_L\}}
\] (5.16)
\[
F_j := \max_{i \in \{\tau_j, ..., \tau_{j-1}\}} \left| \sum_{i=\tau_j}^{\tau_{j-1}} (-1)^{i-1} V_i + (-1)^{i-1} \frac{V_j}{2} \right| = \max_{i \in \{0, ..., \tau_j-1\}} \left| \sum_{i=0}^{\tau_j-1} (-1)^{i-1} V_i + (-1)^{i-1} \frac{V_j}{2} \right|
\]

In this step we prove (5.15) subject to Claims 5.1 and 5.2 and to Lemma 5.3 below.
Claim 5.1. For every $c > 0$ and $\eta > 0$,

$$\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta, \mu_p} \left( \max_{j \leq L^{1/3}} |T_{j, \varepsilon}| \geq \eta L^{1/3} \right) = 0. \quad (5.17)$$

Claim 5.2. For every $c > 0$ and $\eta > 0$,

$$\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta, \mu_p} \left( \max_{j \leq L^{1/3}} F_j \geq \eta L^{1/3} \right) = 0. \quad (5.18)$$

Lemma 5.3 (Lemma 4.12 Carmona and Pétrélis (2016)). For $\beta > 0$ and $\varepsilon > 0$ there exists a $c_\varepsilon > 0$ such that for $L \in \mathbb{N}$,

$$P_{\beta, \mu_p}(v_L \geq c_\varepsilon L^{1/3}) \leq \varepsilon.$$ 

Lemma 5.3 is the same as Lemma 4.12 in Carmona and Pétrélis (2016) except that it is stated there under $P_\beta$ instead of $P_{\beta, \mu_p}$ but the proof is literally the same and we will not repeat it here.

We first recall (5.16) and we observe that when $L \in \mathcal{X}$ then $\xi_L \equiv \tau_{v_L} - 1$ and we can write

$$\max_{n \leq v_L} |M_n - M_{L/k}(n)| = \max_{n \leq v_L-1} |M_n - M_{L/k}(n)| \leq \max_{j \leq v_L} |T_{j, \varepsilon}| + \max_{j \leq v_L, X_j \leq \varepsilon} F_j. \quad (5.19)$$

Therefore, (5.15) (and consequently Proposition 2.4) will be proven once we show that for every $\eta > 0$

$$\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta, \mu_p} \left( \max_{j \leq v_L} |T_{j, \varepsilon}| \geq \eta L^{1/3} \right) = 0, \quad (5.20)$$

and also

$$\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta, \mu_p} \left( \max_{j \leq v_L} F_j \geq \eta L^{1/3} \right) = 0. \quad (5.21)$$

We recall that

$$v_L = \max \{ j \geq 1 : X_1 + \cdots + X_j \leq L \}, \quad (5.22)$$

and we set

$$v'_L \equiv v_L - \min \{ j \geq 1 : X_1 + \cdots + X_j \geq \frac{L}{2} \}. \quad (5.23)$$

We observe that

$$\max_{j \leq v_L} |T_{j, \varepsilon}| \leq \max_{j \leq v'_L} |T_{j, \varepsilon}| + \max_{j \in \{ v_L+1, \ldots, v'_L \}} |T_{j, \varepsilon}| \quad (5.24)$$

but we can bound from above

$$\max_{j \in \{ v_L+1, \ldots, v'_L \}} |T_{j, \varepsilon}| \leq |T_{v'_L, \varepsilon}| + \max_{j \in \{ v_L+1, \ldots, v'_L \}} \left( \sum_{r = v_L+1}^{v'_L} M_{\text{exc}}(r) 1_{\{ X_r \leq \varepsilon \}} \right) \quad (5.25)$$

$$\leq |T_{v'_L, \varepsilon}| + \sum_{r = v_L+1}^{v'_L} M_{\text{exc}}(r) 1_{\{ X_r \leq \varepsilon \}} + \max_{j \in \{ v_L+2, \ldots, v'_L \}} \left( \sum_{r = j}^{v'_L} M_{\text{exc}}(r) 1_{\{ X_r \leq \varepsilon \}} \right).$$

At this stage, we note that

$$v_L + 1 - v'_L = 1 + \min \{ j \geq 1 : X_1 + \cdots + X_j \geq \frac{L}{4} \} \quad (5.26)$$

and then either $\varepsilon L/4 \in \mathcal{X}$ and the r.h.s. in (5.26) equals $v_L/4 + 1$ or the r.h.s. in (5.26) equals $v_L/4 + 2$. In this last case, we note that $v_L/4 + 2 \leq v_L/2 + 1$ except if $v_L/4 = v_L/2$ but this means that $X_{v_L/4+1} = X_{v_L/2+1} > L/4$
and since the excursions associated with a geometric area larger than $L/k$ are not taken into account in the present computation it suffices to choose $k \geq 5$ to make sure that

\[
\max_{j \in \{v_{1/2} + 1, \ldots, v_1\}} \left| \sum_{r \in j} M^{\text{exc}}(r)1_{\{X_r \leq t \}} \right| \leq \max_{j \in \{1, \ldots, v_{3L/4}\}} \left| \sum_{r \in j} M^{\text{exc}}(r + 1 - r)1_{\{X_{r+1} \leq t \}} \right|. \tag{5.27}
\]

We can finally use (5.24–5.27) to conclude that

\[
\max_{j \leq v_1} \left| T_{j, \frac{1}{2}} \right| \leq \max_{j \geq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| + 2 \max_{j \in \{1, \ldots, v_{3L/4}\}} \left| \sum_{r \in j} M^{\text{exc}}(r + 1 - r)1_{\{X_{r+1} \leq t \}} \right|. \tag{5.28}
\]

In the same spirit we bound from above

\[
\max_{j \leq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| = \text{law} \max_{j \in \{1, \ldots, v_{3L/4}\}} \left| \sum_{r \in j} M^{\text{exc}}(r + 1 - r)1_{\{X_{r+1} \leq t \}} \right| \leq \max_{j \geq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| = \text{law} \max_{j \geq v_{3L/4}} \left| F_{j+1} \right|. \tag{5.30}
\]

By reversibility, we note that, under $P_{\beta, \mu}$ ($|L| \in \mathcal{X}$), the following equalities in distribution hold true

\[
\max_{j \leq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| = \text{law} \max_{j \in \{1, \ldots, v_{3L/4}\}} \left| \sum_{r \in j} M^{\text{exc}}(r + 1 - r)1_{\{X_{r+1} \leq t \}} \right| = \max_{j \geq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| = \text{law} \max_{j \geq v_{3L/4}} \left| F_{j+1} \right|. \tag{5.30}
\]

Thus, we can conclude from (5.28–5.30) that (5.20) and (5.21) will be proven once we show that for every $\eta > 0$ we have

\[
\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta, \mu} \left( \max_{j \leq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| \geq \eta L^{1/3} \left| L \in \mathcal{X} \right. \right) = 0,
\]

\[
\lim_{k \to \infty} \limsup_{L \to \infty} P_{\beta, \mu} \left( \max_{j \leq v_{3L/4}} \left| T_{j, \frac{1}{2}} \right| \geq \eta L^{1/3} \left| L \in \mathcal{X} \right. \right) = 0. \tag{5.31}
\]

At this stage, the 2 inequalities in (5.31) are straightforward consequences of Lemma 5.3 above, of Lemma 5.4 and of Claims 5.1 and 5.2. Lemmas 5.4 and 5.3 indeed imply that it is sufficient to prove both inequalities in (5.31) without the conditioning $\{L \in \mathcal{X}\}$ and with $\epsilon L^{1/3}$ instead of $v_{3L/4}$ so that we are left with Claims 5.1 and 5.2.

**Lemma 5.4.** For every $\beta > 0$, there exists a $M > 0$ such that, for every function $G : \cup_{k=1}^{\infty} \mathbb{R}_k^+ \to \mathbb{R}_+$ and every $L \in \mathbb{N}$, we have

\[
E_{\beta, \mu} \left[ G(V_0, \ldots, V_{v_{3L/4} - 1}) \left| L \in \mathcal{X} \right. \right] \leq M E_{\beta, \mu} \left[ G(V_0, \ldots, V_{v_{3L/4} - 1}) \right]. \tag{5.32}
\]

**Proof.** We compute the Radon Nikodym density of the image measure of $P_{\beta, \mu} \left( |L| \in \mathcal{X} \right)$ by $(V_0, \ldots, V_{v_{3L/4} - 1})$ w.r.t. its counterpart without conditioning. For $y \in \{1, \ldots, 2/3\}, t \in \{1, \ldots, y\}, m \in \{0, \ldots, t\}$ and $(z_0, z_1, \ldots, z_{t-1}) \in \mathbb{Z}_t$ satisfying $t + |z_0| + \cdots + |z_{t-1}| = y$ we obtain

\[
P_{\beta, \mu} \left( v_{3L/4} = m, S_m = y, \tau_m = t, (V_0, \ldots, V_{v_{3L/4} - 1}) = (z_0, \ldots, z_{t-1}) \right| L \in \mathcal{X} \right) = G_L(y) + K_L(y),
\]

with

\[
G_L(y) = \sum_{n=0}^{L/8} P_{\beta, \mu} \left( n \in \mathcal{X} \right) P_{\beta, \mu} \left( X = L - n - y \right),
\]

\[
K_L(y) = \sum_{n=1+L/8}^{L/4} P_{\beta, \mu} \left( n \in \mathcal{X} \right) P_{\beta, \mu} \left( X = L - n - y \right) P_{\beta, \mu} \left( X \geq \frac{3L}{4} - y \right). \tag{5.33}
\]
The rest of the proof consists in showing that $G_n(y)$ and $K_n(y)$ are bounded above uniformly in $L \in \mathbb{N}$ and $y \in \{0, \ldots, 3L/4\}$. We will focus on $G_n$ since $K_n$ can be treated similarly. The constants $c_1, \ldots, c_4$ below are positive and independent of $L, n, y$. By recalling (3.2) and since $L - n - y \geq L/4$ when $n \in \{0, \ldots, L/8\}$, we can claim that in the numerator of $G_n(y)$, the term $P_{\mu, \nu}(X = L - n - y)$ is bounded above by $c_1/L^{4/3}$ independently of $n$ while $\sum_{i=0}^{L/4} P_{\mu, \nu}(n \in \mathbb{N}) \leq c_2 L^{1/3}$ for every $L \in \mathbb{N}$. For the denominator, (3.2) tells us that $P_{\mu, \nu}(L \in \mathbb{X}) \geq c_3/L^{2/3}$ and that

$$P_{\mu, \nu}(X \geq \frac{2L}{3} - y) \geq P_{\mu, \nu}(X \geq \frac{2L}{3}) \geq \frac{c_4}{L^{2/3}}, \quad L \in \mathbb{N}, \ y \in \{0, \ldots, 3L/4\}.$$ 

This terminates the proof.

$\square$

Step 2: proof of Claim 5.1

We recall (2.22) and in the Definition 3.2. For every $j \geq 1$ we set

$$R_j := \sum_{i=0}^{\tau_j - \tau_{j-1} - 1} (-1)^{i-1} |V_{\tau_{j-1} + i}| = \sum_{i=0}^{\tau_j - 1} (-1)^{i-1} |V_i|$$

(5.34)

so that $(R_i)_{i \geq 1}$ is an i.i.d. sequence of random variables satisfying $|R_j| = |M^\text{exc}(j)|$ for every $j \geq 1$. We note that for $n \geq 2$,

$$T_n \triangleq T_{n-1, j} + (-1)^{\tau_{n-1}} \{1 \{V_{\tau_{n-1}} = 0\} \epsilon_n + 1 \{V_{\tau_{n-1}} \neq 0\} (\text{sign}(V_{\tau_{n-1}}))\} R_n 1 \{X_{j-1} \leq \frac{k}{n}\}$$

(5.35)

and we define the filtration $(\mathcal{F}_n)_{n \geq 1}$ by

$$\mathcal{F}_n := \sigma(\epsilon_1, \ldots, \epsilon_n, V_{\tau_{j-1} \leq 1}), \quad n \in \mathbb{N}.$$ 

The expectation of $T_{n, k}$ conditioned by $\mathcal{F}_{n-1}$ is easily computed since $T_{n-1, j}$ is $\mathcal{F}_{n-1}$ measurable. Therefore, for $n \geq 2$ we obtain

$$E_{\mu, \nu} \left[ T_{n, k} | \mathcal{F}_{n-1} \right] = T_{n-1, j} - (-1)^{\tau_{n-1}} \text{sign}(V_{\tau_{n-1}}) E_{\mu, \nu} \left[ 1 \{V_{\tau_{n-1}} \neq 0\} R_n 1 \{X_{j-1} \leq \frac{k}{n}\} \right],$$

(5.36)

and with the subsequent notation

$$\varphi_1(x) = E_{\mu, \nu} \left[ 1 \{V_{\tau_{n-1}} \neq 0\} R_n 1 \{X_{j-1} \leq \frac{k}{n}\} \right], \quad x \in \mathbb{N}.$$ 

(5.37)

we can rewrite $T_{n, k} = Q_{n, k} - J_{n, k}$ such that

$$Q_{n, k} := M^\text{exc}(1) \{X_{j-1} \leq \frac{k}{n}\} + \sum_{r=2}^{n} M^\text{exc}(r) \{X_{j-1} \leq \frac{k}{n}\} (-1)^{\tau_{n-1}} \text{sign}(V_{\tau_{n-1}}) \varphi_1(L/k)$$

(5.38)

and

$$J_{n, k} := \varphi_1(L/k) \sum_{r=2}^{n} (-1)^{\tau_{n-1}} \text{sign}(V_{\tau_{n-1}}).$$

(5.39)

Equation 5.35 guaranties that $(Q_{n, k})_{n \in \mathbb{N}}$ is an $L^2$ martingale. The proof of Claim 5.1 will therefore be proven once we show that

$$\lim_{k \to \infty} \lim_{L \to \infty} \max_{n \in \mathbb{N}} \left| Q_{n, k} \right| \geq \eta L^{1/3} = 0,$$ 

(5.40)
so that (5.45) becomes
\[ \lim_{k \to \infty} \limsup_{\beta \to \infty} \mathbb{P}_{\beta, \mu_\beta} \left( \max_{n \leq L \log \beta} |J_{n, \beta}^i| \geq \eta L^{1/3} \right) = 0. \] (5.41)

We begin with (5.40) and we apply Doob inequality with the fact that \((Q_{n, \beta})_{n \in \mathbb{N}}\) is an \(L^2\) martingale to assert that there exists \(c_1 > 0\) such that
\[ \mathbb{P}_{\beta, \mu_\beta} \left( \max_{n \leq L \log \beta} |Q_{n, \beta}^i| \geq \eta L^{1/3} \right) \leq \frac{4 \mathbb{E}_{\beta, \mu_\beta} \left[ Q_{n, \beta}^2 \right]}{\eta^2 L^{2/3}} \] (5.42)
and that
\[ \mathbb{E}_{\beta, \mu_\beta} \left[ Q_{n, \beta}^2 \right] = \mathbb{E}_{\beta, \mu_\beta} \left[ (M^{\text{exc}}(1))^2 1_{\{X_i, r \leq \beta\}} \right] + \sum_{r=2}^{L^{1/3}} \mathbb{E}_{\beta, \mu_\beta} \left[ (M^{\text{exc}}(r)) 1_{\{X_i, r \leq \beta\}} \right] \] (5.43)

At this stage, we recall (2.22) and (5.34) which yield that \((M^{\text{exc}}(r))^2 = R_r^2\) for every \(r \geq 1\). Moreover \((R_r)_{r \geq 1}\) is an i.i.d. sequence of random variables and therefore, we deduce from (5.43) that
\[ \mathbb{E}_{\beta, \mu_\beta} \left[ Q_{n, \beta}^2 \right] \leq 2cL^{1/3} \left( \mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right] + \varphi_1(L/k)^2 \right) \leq 4cL^{1/3} \mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right] \] (5.44)
where the second inequality above is the result of Jensen inequality. As a result, we need to bound from above the quality \(\mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right]\), i.e.,
\[ \mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right] = \mathbb{E}_{\beta, \mu_\beta} \left[ \left( -\frac{V_0}{2} + \sum_{i=1}^{\tau_i-1} (-1)^{i-1} \frac{V_i - V_{i-1}}{2} + (-1)^{i-2} \frac{V_{\tau_i-1}}{2} \right)^2 1_{\{X_i, r \leq \beta\}} \right]. \] (5.45)

At this stage, we substitute an expectation with respect to \(\mathbb{P}_\beta\) to that w.r.t. \(\mathbb{P}_{\beta, \mu_\beta}\) in the r.h.s. of (5.45). We proceed as follow. We define, for \(y \in \mathbb{N}\), the set of excursions \(\mathcal{G}_y := \sup_{i \geq 1} \mathcal{G}_{i, y}\), defined with
\[ \mathcal{G}_{i, y} := \cup_{j=1}^{i-1} \{ (v_i)_{i=0}^{-1} : v_i = 0 \ \forall \ i \leq j - 1 \ \text{and} \ v_i > 0 \ \forall \ i \geq j \ \text{and} \ s + \sum_{i=0}^{r-1} |v_i| \leq x \} \] (5.46)
so that (5.45) becomes
\[ \mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right] = 2 \sum_{i=1}^{\tau_i-1} \sum_{y \in \mathcal{G}_{i, y}} \mathbb{P}_{\beta, \mu_\beta} \left[ (V_i)_{i=0}^{-1} = (v_i)_{i=0}^{-1} \ \text{and} \ V_i \leq 0 \right] \] (5.47)
\[ \times \left( -\frac{V_0}{2} + \sum_{i=1}^{\tau_i-1} (-1)^{i-1} \frac{V_i - V_{i-1}}{2} + (-1)^{i-2} \frac{V_{\tau_i-1}}{2} \right)^2, \]
where the factor 2 in front of the r.h.s. comes from the fact that negative and positive excursions contribute the same when computing \(\mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right]\). At this stage, we recall (3.1) and we observe that for all \(v \in \mathcal{G}_{i, y}\) we have
\[ \mathbb{P}_{\beta} \left[ (V_i)_{i=0}^{-1} = (v_i)_{i=0}^{-1} \ \text{and} \ V_i = 0 \right] = \frac{1}{c_2} \left( 1_{\{v_0 = 0\}} + \frac{1}{2} 1_{\{v_0 \neq 0\}} \right) \mathbb{P}_{\beta, \mu_\beta} \left[ (V_i)_{i=0}^{-1} = (v_i)_{i=0}^{-1} \ \text{and} \ V_i \leq 0 \right]. \] (5.48)
It remains to combine (5.47) with (5.48) to obtain that there exit \(c_2 > 0\) such that
\[ \mathbb{E}_{\beta, \mu_\beta} \left[ R_1^2 1_{\{X_i, r \leq \beta\}} \right] \leq c_2 \mathbb{E}_{\beta} \left[ \left( -\frac{V_0}{2} + \sum_{i=1}^{\tau_i-1} (-1)^{i-1} \frac{V_i - V_{i-1}}{2} + (-1)^{i-2} \frac{V_{\tau_i-1}}{2} \right)^2 1_{\{V_{\tau_i-1} > 0, v_i = 0\}} \right] \left( X_i, r \leq \beta \right) \] (5.49)
For $x \in \mathbb{N}$, we decompose $\varphi_2(x)$ with respect to the value taken by $\tau_1$, i.e.,

$$
\varphi_2(x) = \sum_{s=2}^x E_\rho\left[ \left( \sum_{i=1}^s (-1)^i U_i \right)^2 1_{\{V_i, v_i = v = 0, V_i = 0\}} 1_{\{s, v_i = 0\}} 1_{\{x_1 \leq x\}} \right] := \sum_{s=2}^x \alpha_{s,x}
$$

(5.50)

We define for $j \in \mathbb{N}$ and $V \in \{0\} \times 2^\mathbb{N}$ the geometric area seen from the minimum of $V$ after $j$ steps, i.e.,

$$
A_{\min,j}(V) = j + \sum_{i=1}^j |V_i - \min\{0, V_1, \ldots, V_j\}|
$$

(5.51)

and we define also

$$
\mathcal{U}_{s,x} := \{\{U_i\}_{i=1}^s : V_i - 0, \tau_1 = s, V_s = 0, X_i \leq x\}
$$

$$
\mathcal{O}_{s,x} := \{\{U_i\}_{i=1}^s : V_s = 0, A_{\min,s}(V) \leq x\}
$$

(5.52)

and we apply to $\mathcal{U}_{s,x}$ the $s$ shifts $\theta^j, j \in \{0, \ldots, s-1\}$ defined by

$$
\theta^j(u_1, \ldots, u_s) = (u_{j+1}, \ldots, u_s, u_1, \ldots, u_j)
$$

The crucial point here is that for every $(u_1, \ldots, u_s) \in \mathcal{U}_{s,x}$ and every $0 \leq j \leq s-1$,

(a) $\theta^j(u_1, \ldots, u_s) \in \mathcal{O}_{s,x}$

(b) $P_\rho(\mathcal{U}_{s,x}) = (u_1, \ldots, u_s)) = P_\rho(\mathcal{U}_{s,x}) = \theta^j(u_1, \ldots, u_s))

(c) $A_{\min,s}(u_1, \ldots, u_s) = A_{\min,s}(\theta^j(u_1, \ldots, u_s))

(d) If $s \in 2\mathbb{N}$ for $(v_1, \ldots, v_s) = \theta^j(u_1, \ldots, u_s)$ we have $\left( \sum_{i=1}^s (-1)^i v_i \right)^2 = \left( \sum_{i=1}^s (-1)^i u_i \right)^2

(e) For every $(v_1, \ldots, v_s) \in \mathcal{U}_{s,x} \theta^j(\mathcal{U}_{s,x})$ there exists a unique $0 \leq j \leq s-1$ and a unique $(u_1, \ldots, u_s) \in \mathcal{U}_{s,x}$ such that $(v_1, \ldots, v_s) = \theta^j(u_1, \ldots, u_s).

As a consequence, for every $s \in 2\mathbb{N}$ and $x \in \mathbb{N}$ we have the upper bound

$$
\alpha_{s,x} \leq \frac{1}{s} E_\rho\left[ \left( \sum_{i=1}^s (-1)^i U_i \right)^2 1_{\{V_i = 0\}} 1_{\{A_{\min,s}(V) \leq s\}} \right],
$$

(5.53)

and moreover, one can show that there exists a $c_3 > 0$ such that for $s \in 2\mathbb{N} + 1$ and $x \in \mathbb{N}$ we have

$$
\alpha_{s,x} \leq c_3 \alpha_{s+1,x+1},
$$

(5.54)

since it suffices to add one increment equal to 0 in front of a trajectory from $\mathcal{U}_{s,x}$ to obtain a trajectory from $\mathcal{U}_{s+1,x+1}$. We recall (5.50) and, as a consequence of (5.53) and (5.54), we can claim that

$$
\varphi_2(x) \leq \psi(x) + c_3 \psi(x + 1), \quad x \in \mathbb{N},
$$

(5.55)

with

$$
\psi(x) := \sum_{p \geq 1} \frac{1}{2p} E_\rho\left[ \left( \sum_{i=1}^{2p} (-1)^i U_{i} \right)^2 1_{\{V_{2p} = 0\}} 1_{\{A_{\min,2p}(V) \leq x\}} \right] := \sum_{p \geq 1} \frac{1}{2p} \gamma_{2p,x}.
$$

(5.56)

We easily conclude that

$$
\gamma_{2p,x} \leq 2 E_\rho\left[ \left( \sum_{i=1}^{p} (-1)^i U_{i} \right)^2 1_{\{V_{2p} = 0\}} 1_{\{A_{\min,2p}(V) \leq x\}} \right] + 2 E_\rho\left[ \left( \sum_{i=p+1}^{2p} (-1)^i U_{i} \right)^2 1_{\{V_{2p} = 0\}} 1_{\{A_{\min,2p}(V) \leq x\}} \right]
$$

$$
= 4 E_\rho\left[ \left( \sum_{i=1}^{p} (-1)^i U_{i} \right)^2 1_{\{V_{2p} = 0\}} 1_{\{A_{\min,2p}(V) \leq x\}} \right] \leq 4 E_\rho\left[ \left( \sum_{i=1}^{p} (-1)^i U_{i} \right)^2 1_{\{V_{2p} = 0\}} 1_{\{A_{\min}(V) \leq x\}} \right]
$$

$$
:= 4 A_{p,x},
$$

(5.57)
Thus, (5.59–5.62) allow us to assert that there exists a stopping time with respect to \( P_p \) that

\[
A_{p,x} = \sum_{y \leq x} E_p \left[ \left( \sum_{i=1}^{n} (-1)^i U_i \right)^2 1_{\{V_{i} = y\} 1_{\{A_{\max}(V) \leq x\}} \} \right] P_p (V_p = y)
\]  

(5.58)

and it remains to use a local central limit theorem in (Durrett, 2010, Theorem 3.5.2) to claim that there exists a \( c_5 > 0 \) such that for every \( y \in \mathbb{Z} \) we have that \( P_p (V_p = y) \leq \frac{c_5}{y^2} \). Finally

\[
\psi(x) \leq c_5 \sum_{p \geq 1} \frac{1}{p^{3/2}} E_p \left[ \left( \sum_{i=1}^{n} (-1)^i U_i \right)^2 1_{\{A_{\max}(V) \leq x\}} \right].
\]  

(5.59)

At this stage, we let \((\mathcal{G}_n)_{n \geq 0}\) be the natural filtration associated with \((U_i)_{i \in \mathbb{N}}\) and we set

\[
\bar{\tau}_x := \inf \{ j \geq 1 : A_{\min,j}(V) \geq x \},
\]

(5.60)

which is a stopping time with respect to \((\mathcal{G}_n)_{n \geq 1}\). For every \( p \in \mathbb{N} \), the inequality \( A_{\min,p}(V) \leq x \) implies that \( p \leq \bar{\tau}_x \) and therefore

\[
E_p \left[ \left( \sum_{i=1}^{n} (-1)^i U_i \right)^2 1_{\{A_{\max}(V) \leq x\}} \right] \leq E_p \left[ \left( \sum_{i=1}^{\bar{\tau}_x} (-1)^i U_i \right)^2 \right].
\]

(5.61)

Using that \( \left( \sum_{i=1}^{n} (-1)^i U_i \right)^2 - nE_p (U_i^2) \) is a \((\mathcal{G}_n)_{n \geq 1}\) martingale, we can assert that for every \( p \in \mathbb{N} \),

\[
E_p \left[ \left( \sum_{i=1}^{\bar{\tau}_x} (-1)^i U_i \right)^2 \right] = E_p (\bar{\tau}_x \wedge p) E_p (U_i^2) - E_p (\bar{\tau}_x) E_p (U_i^2).
\]

(5.62)

Thus, (5.59–5.62) allow us to assert that there exists a \( c_6 > 0 \) and \( c_7 > 0 \) such that

\[
\psi(x) \leq c_6 \sum_{p = 1}^{\infty} \frac{1}{p^{3/2}} E_p (\bar{\tau}_x \wedge p) = c_6 \sum_{p = 1}^{\infty} \frac{1}{p^{3/2}} E_p (p 1_{\bar{\tau}_x \geq p} + \bar{\tau}_x 1_{\bar{\tau}_x < p})
\]

(5.63)

\[
\leq c_6 E_p \left[ \sum_{p = 1}^{\infty} \frac{1}{p^{3/2}} + \bar{\tau}_x \sum_{p = 1}^{\infty} \frac{1}{p^{3/2}} \right] \leq c_7 E_p (\sqrt{\bar{\tau}_x}) \leq c_7 \sqrt{E_p (\bar{\tau}_x)},
\]

so that finally (5.55) and (5.63) yield that there exists \( c_8 > 0 \) such that for every \( x \in \mathbb{N} \),

\[
\varphi_2 (x) \leq c_8 \sqrt{E_p (\bar{\tau}_x)}, \quad x \in \mathbb{N}.
\]

(5.64)

At this stage, we combine (5.42), (5.44), (5.49) with (5.64) (at \( x = L/k \)) and we obtain that there exists \( c_9 > 0 \) such that

\[
P_{p,\mu_y} \left( \max_{n \leq L/k} |Q_n| \geq \eta L^{1/3} \right) \leq c_9 \frac{\sqrt{E_p (\bar{\tau}_x)}}{\eta L^{1/3}}.
\]

Thus, we complete the proof of (5.40) with a straightforward application of Lemma 5.5 (proven in Step 5).

We continue with the proof of (5.41). We apply Cauchy Schwartz to (5.37) and we recall (5.49) to conclude that there exists a \( c_1 > 0 \) such that

\[
\varphi_1(x) \leq E_{p,\mu_y} \left[ R_1^2 1_{\{X_1 \leq x\}} \right]^{1/2} \leq c_1 \varphi_2 (x)^{1/2}.
\]

(5.65)
Then, we use (5.64) and Lemma 5.5 to conclude that there exists a $c_2 > 0$ such that

$$ \varphi_1(x) \leq cx^{1/6}, \quad x \in \mathbb{N}. $$

(5.66)

We recall

$$ J_{n,k} := \varphi_1(L/k) \sum_{r=2}^{n} (-1)^{\tau_{r-1}} \text{sign}(V_{\tau_{r-1}}) $$

and therefore, (5.41) will be proven once we show that

$$ \left( \frac{1}{\sqrt{L/\epsilon}} \max_{n \leq L/\epsilon} \left| \sum_{r=2}^{n} (-1)^{\tau_{r-1}} \text{sign}(V_{\tau_{r-1}}) \right| \right)_{L \in \mathbb{N}} $$

is a tight sequence of random variables.

The idea to perform this proof consists in rewriting the sum in (5.67) as a sum of i.i.d. centered random variables with a finite second moment. Thus, the tightness of the sequence of random variables in (5.67) is a straightforward consequence of Donsker invariance principle.

Step 3: proof of Claim 5.2

We set

$$ B_{L, \eta} := \left\{ \max_{j \leq L/\eta, X_j \leq 1/\epsilon} F_j < \eta L^{1/3} \right\} \cap \left\{ \{X_j > 1/\epsilon\} \cup \{F_j < \eta L^{1/3}, X_j \leq 1/\epsilon\} \right\}, $$

(5.70)
and we use (5.16) to recall that the sequence \((F_i)_{i \in \mathbb{N}}\) is i.i.d. so that
\[
\mathbb{P}_{\beta, \mu \eta}(B_{L, \eta}) = e^{L^{1/3} \log \left(1 - \mathbb{P}_{\mu \eta}(F_{i} \geq \eta L^{1/3}, x_{i} \leq \frac{1}{\kappa})\right)}.
\]
(5.71)

Thus, (5.71) guarantees that the proof of Claim 5.2 will be complete once we show that for every \(\eta > 0\),
\[
\lim_{k \to \infty} \limsup_{L \to \infty} L^{1/3} \mathbb{P}_{\beta, \mu \eta}(F_{i} \geq \eta L^{1/3}, x_{i} \leq \frac{1}{\kappa}) = 0
\]
(5.72)

By using the mapping of trajectories introduced in (5.45–5.49) we again substitute the law \(\mathbb{P}_{\beta}\) to \(\mathbb{P}_{\beta, \mu \eta}\) in the r.h.s. of (5.72). We indeed obtain that there exists a \(c > 0\) such that
\[
\mathbb{P}_{\beta, \mu \eta}(F_{i} \geq \eta L^{1/3}, x_{i} \leq \frac{1}{\kappa}) \leq c \mathbb{P}_{\beta}(C_{L, \eta})
\]
(5.73)

where
\[
C_{L, \eta} := \left\{ F_{i} \geq \eta L^{1/3}, V_{\tau_{i}-1} > 0, V_{\tau_{i}} = 0, x_{i} \leq 1 + \frac{1}{\kappa} \right\}
\]
(5.74)

and with an alternative description of \(F_{i}\), i.e.,
\[
F_{i} := \frac{1}{2} \max_{j \in \{1, \ldots, \tau_{i}\}} \left| \sum_{s=1}^{j} (-1)^{s-1} U_{i, s} \right|.
\]
(5.75)

We slightly modify the notations in (5.52), i.e., for \(x \in \mathbb{N}\),
\[
\overline{\mathcal{W}}_{s, x} := \{(U)_{i=1}^{s} : V_{s-1} > 0, \tau_{s} = s, V_{s} = 0, x_{1} \leq x, F_{i} \geq \eta L^{1/3}\}
\]
(5.76)

and we note that
\[
C_{L, \eta} := \bigcup_{s \geq 2} \overline{\mathcal{W}}_{s, \frac{L}{\kappa}}.
\]
(5.77)

We apply to \(\mathcal{W}_{s, x}\) the \(s-1\) shifts \(\theta^{j}, j \in \{0, \ldots, s-1\}\) defined by
\[
\theta^{j}(u_{1}, \ldots, u_{s}) = (u_{j+1}, \ldots, u_{s}, u_{1}, \ldots, u_{j})
\]

The crucial point here is that for every \((u_{1}, \ldots, u_{s}) \in \overline{\mathcal{W}}_{s, x}\) and every \(0 \leq j \leq s-1\), the properties (a–c) and (e) stated below (5.52) are still satisfied here with \(\mathcal{W}_{s, x}\) and \(\overline{\mathcal{W}}_{s, x}\) whereas the (d) property is replaced by
\[
\max_{j \in \{1, \ldots, s\}} \left| \sum_{i=1}^{j} (-1)^{j-1} V_{i} \right| \geq \frac{1}{3} \max_{j \in \{1, \ldots, s\}} \left| \sum_{i=1}^{j} (-1)^{j-1} U_{i} \right|,
\]
(5.78)

with \((v_{1}, \ldots, v_{s}) = \theta^{j}(u_{1}, \ldots, u_{s})\). As a consequence, we obtain the following upper bound,
\[
\mathbb{P}_{\beta}(C_{L, \eta}) \leq \sum_{s \geq 2} \frac{1}{s} \mathbb{P}_{\beta}(F_{i} \geq \frac{\eta L^{1/3}}{3}, V_{s} = 0, A_{\min, s} \leq \frac{L}{\kappa}).
\]
(5.79)

At this stage, we consider a sequence of \(s\) increments \((U)_{i=1}^{s}\) such that the associated \(V\) trajectory satisfies \(V_{s} = 0\) and \(A_{\min, s}(V) \leq \frac{L}{\kappa}\). Then, the \(V\) trajectory defined by \(V_{i} = V_{s-i}\) for \(i \in \{0, \ldots, s\}\) has increments \((-U_{s-i})_{i=1}^{s}\) also satisfies \(V_{s} = 0\) and \(A_{\min, s}(V) \leq \frac{L}{\kappa}\). We can use this auxiliary trajectory and check easily that \(F_{i} \leq \max\{F_{1,1}, F_{1,2}\}\) where
\[
F_{1,1} := \max_{j \in \{1, \ldots, |s|+1\}} \left| \sum_{i=1}^{j} (-1)^{j-1} U_{i} \right| \quad \text{and} \quad F_{1,2} := \max_{j \in \{1, \ldots, |s|+1\}} \left| \sum_{i=1}^{j} (-1)^{j-1} U_{i} \right|.
\]
(5.80)
Moreover,
\[
F_{1,2} \leq \max_{j \in \{-\frac{1}{3}+1, \ldots, \frac{1}{3}+1\}} \left| \sum_{i=1}^{s} (-1)^{i-1} U_i + (-1)^{s+1} \sum_{i=1}^{s} (-1)^{i-1} (-U_{s+1-i}) \right| \leq \left| \sum_{i=1}^{s} (-1)^{i-1} U_i \right| + \max_{j \in \{-\frac{1}{3}+1, \ldots, \frac{1}{3}+1\}} \left| \sum_{i=1}^{s} (-1)^{i-1} (-U_{s+1-i}) \right| := \sum_{i=1}^{s} (-1)^{i-1} U_i + \bar{F}_{1,1},
\]
(5.81)
and a straightforward computations gives us that \( \sum_{i=1}^{s} (-1)^{i-1} U_i \leq F_{1,1} + \bar{F}_{1,1} \). Thus,
\[
F_{1} \leq 3 \max\{F_{1,1}, \bar{F}_{1,1}\}
\]
(5.82)
and we note that, conditioned on \( \bar{V} \equiv 0 \) and \( A_{\min,t}(\bar{V}) \leq \frac{1}{2} \), the two random variables \( F_{1,1} \) and \( \bar{F}_{1,1} \) have the same law. As a consequence, \( F_{1} \) can be replaced by \( \bar{F}_{1,1} \) in the r.h.s of (5.79) and the proof will be complete once we prove that for every \( \eta > 0 \),
\[
\lim_{k \to \infty} \limsup_{L \to \infty} L^{1/3} \sum_{s \geq 2} \frac{1}{s} \alpha_{\eta,s,L,k} = 0
\]
(5.83)
where
\[
\alpha_{\eta,s,L,k} := P_{\beta}(F_{1,1} > \eta L^{1/3}, V_k = 0, A_{\min,t} \leq \frac{L}{k}).
\]
By Markov inequality applied at time \( t_k := [\frac{k}{2}] + 1 \) we can write that for every \( s \geq 2 \),
\[
\alpha_{\eta,s,L,k} \leq \sum_{y \in \mathbb{Z}} P_{\beta}(F_{1,1} > \eta L^{1/3}, V_{t_k} = y, A_{\min,t_k} \leq \frac{L}{k}) P_{\beta}(V_{t_k} = y)
\]
(5.84)
so that it remains to use the local central limit Theorem in (Durrett, 2010, Theorem 3.5.2) to claim that there exists a \( c > 0 \) such that for every \( y \in \mathbb{Z} \) we have that \( P_{\beta}(V_{t_k} = y) \leq \frac{c}{\sqrt{2}} \) and to sum over \( y \) to obtain
\[
\sum_{s \geq 2} \frac{1}{s} \alpha_{\eta,s,L,k} \leq \frac{c}{\sqrt{2}} \sum_{s \geq 2} \frac{1}{s^{3/2}} P_{\beta}(F_{1,1} > \eta L^{1/3}, A_{\min,t_k} \leq \frac{L}{k})
\]
\[
\leq 2c \sum_{p \geq 1} \frac{1}{p^{3/2}} P_{\beta}\left( \max_{j \in \{1, \ldots, p+1\}} \left| \sum_{i=1}^{j} (-1)^{i-1} U_i \right| > \eta L^{1/3}, A_{\min,p+1} \leq \frac{L}{k} \right)
\]
(5.85)
\[
\leq 2c \sum_{p \geq 1} \frac{1}{p^{3/2}} P_{\beta}\left( \max_{j \in \{1, \ldots, p\}} \left| \sum_{i=1}^{j} (-1)^{i-1} U_i \right| > \eta L^{1/3}, A_{\min,p} \leq \frac{L}{k} \right)
\]
(5.86)
where the second inequality in (5.85) is obtained by noting that \( t_{2p} = t_{2p+1} \) for every \( p \geq 1 \). At this stage, we recall the definition of \( \bar{\tau} \) in (5.60) and we recall also that for every \( p \in \mathbb{N} \), the inequality \( A_{\min,p}(V) \leq \frac{L}{k} \) implies that \( p \leq \bar{\tau}_{L/k} \) and therefore
\[
\left\{ \max_{j \in \{1, \ldots, p\}} \left| \sum_{i=1}^{j} (-1)^{i-1} U_i \right| > \eta L^{1/3}, A_{\min,p} \leq \frac{L}{k} \right\} \subset \left\{ \max_{j \in \{1, \ldots, p\}} \sum_{i=1}^{j} (-1)^{i-1} U_i \right\} > \eta L^{1/3} \right\}
\]
(5.87)
Moreover \( \bar{\tau}_{L/k} \) is a stopping time and \( \left( \sum_{i=1}^{n} (-1)^{i-1} U_i \right)_{n \geq 1} \) is a martingale so that by Doob inequality we can claim that
\[
P_{\beta}\left( \max_{j \in \{1, \ldots, p\}} \left| \sum_{i=1}^{j} (-1)^{i-1} U_i \right| > \eta L^{1/3} \right) \leq \frac{E_{\beta}\left( \sum_{i=1}^{p} (-1)^{i-1} U_i \right)^2}{\eta^2 L^{2/3}} = \frac{E_{\beta}(\bar{\tau}_{L/k} \wedge p) E_{\beta}(U_1^2)}{\eta^2 L^{2/3}}
\]
(5.88)
where we have used that \( \sum_{i=1}^{n} (-1)^{i-1} U_i \) \( \to -n E_p(U^2) \) is a martingale. At this stage, (5.83) becomes

\[
\lim_{k \to \infty} \limsup_{L \to \infty} L^{1/3} \sum_{p \geq \beta} \frac{1}{P^{3/2}} \frac{E_p(\bar{T}_{L/k} \wedge p)}{\eta^2 L^{2/3}} = 0, \tag{5.89}
\]

so that, by mimicking (5.63), it remains to prove that

\[
\lim_{k \to \infty} \limsup_{L \to \infty} L^{1/3} \sqrt{\bar{T}_{L/k}} = 0, \tag{5.90}
\]

which is a straightforward consequence of Lemma 5.5 proven in Step 5 below.

**Step 4: Lemma 5.5**

In this section, we state and prove a lemma that allows us to control the growth of \( \bar{T}_x \) as \( x \to \infty \).

**Lemma 5.5.** For every \( \beta > 0 \), there exists a \( c > 0 \) such that \( E_p(\bar{T}_x) \leq cx^{2/3} \) for every \( x \in \mathbb{N} \).

To prove the lemma we need to divide every \( V \) trajectory into pseudo-excursions. To that aim, we define two sequences of random times, i.e.,

\[
\eta := \inf(j \geq \eta_i; V_{j+1} > V_j),
\]

\[
\eta_i := \inf(j > \eta_i; V_j \leq V_{\eta_i}).
\]

The pseudo-excursion indexed by \( j \in \mathbb{N} \) is given by \( (V_{\eta_{i−1}+1}, V_{\eta_{i−1}+2}, \ldots, V_{\eta_i}) \) and we associate it with the quantity

\[
\bar{X}_j = \eta_j - \eta_{j−1} + \sum_{i=\eta_{i−1}+1}^{\eta_{i−1}} V_i - V_{\eta_i}. \tag{5.92}
\]

At this stage, we observe that every pseudo-excursion starts with a non-increasing part of length \( \bar{\eta} \) followed by a real positive excursion (seen from \( V_{\eta} \)) of length \( \eta - \bar{\eta} \). The quantity \( \bar{X} \) corresponds to the total length of the pseudo-excursion plus the area swept by its real excursion. Henceforth, we will abusively call it the area of the pseudo-excursion.

For \( n \in \mathbb{N} \), we denote by \( m_n \) the number of pseudo-excursions that have been completed before time \( n \) and by \( a_n \) the number of pseudo-excursions that have been completed before their cumulated area reaches \( n \), i.e.,

\[
m_n := \max\{j \geq 0; \eta_j \leq n\}, \tag{5.93}
\]

\[
a_n := \max\{j \geq 0; \bar{X}_1 + \cdots + \bar{X}_j \leq n\}. \tag{5.94}
\]

At this stage, we define an increasing functional of the trajectory, i.e.,

\[
R_n = \bar{X}_1 + \cdots + \bar{X}_{m_n} + (n - \eta_{m_n}) + \sum_{i=m_{n−1}+1}^{n} V_i - V_{\eta_{m_n}}. \tag{5.95}
\]

It is easy to see that \( R_n \) is bounded above by \( A_{\min,n} \) for every \( n \in \mathbb{N} \) and every trajectory \( V \) such that \( V_0 = 0 \). Therefore, by recalling the definition of \( \bar{T}_x \) in (5.60) and by defining

\[
\bar{T}_x := \inf\{j \geq 1; R_j(V) \geq x\}, \tag{5.96}
\]

we can claim that \( \bar{T}_x \geq \bar{T}_x \) for every \( x \in \mathbb{N} \) and every \( V \). For this reason, the proof of lemma 5.5 will be complete once we show that there exists a \( c > 0 \) such that

\[
E_p(\bar{T}_x) \leq cx^{2/3}, \quad x \in \mathbb{N}. \tag{5.97}
\]
To prove (5.96), we use the equality 
\[ \frac{1}{x^{2/3}} E_p(\bar{\tau}_x) = E_p\left[ \sum_{j=1}^{a_j} \frac{\eta_j - \eta_{j-1}}{X_j^{2/3}} \frac{\tau_j^{2/3}}{x^{2/3}} \right] + \frac{1}{x^{2/3}} E_p[\bar{\tau}_x - \eta_{a_j}], \] (5.97)

The following claim sheds light on the fact that pseudo-excursions are almost i.i.d.

Claim 5.6. Under \( P_p \) the pseudo-excursions, i.e., \( \{(V_{\eta_1+i} - V_{\eta_{i-1}})_{i=0}^j \} \) are i.i.d. Moreover, for every \( j \geq 1 \) the sequences \( (V_{\eta_1+i} - V_{\eta_{i-1}})_{i=0}^j \) and \( (V_{\bar{\eta}_i+i} - V_{\bar{\eta}_{i-1}})_{i=0}^j \) are independent as well.

Proof. The proof of the claim is a straightforward consequence of strong Markov property combined with the fact that \( 1 + \bar{\eta}_j \) and \( \eta_j \) are stopping times for every \( j \geq 1 \).

Because of Claim 5.6 above, \( (\eta_j - \eta_{j-1})_{j \geq 1} \) and \( (\bar{X}_j)_{j \geq 1} \) are sequences of i.i.d. random variables since, for every \( j \geq 1 \), both \( \eta_j - \eta_{j-1} \) and \( \bar{X}_j \) are functions of the \( j \)-th pseudo excursion only. Thus, by taking the conditional expectation with respect to \( \{a_s, \bar{X}_1, \ldots, \bar{X}_{a} \} \) inside both terms in the r.h.s. of (5.97) we obtain that

\[ \frac{1}{x^{2/3}} E_p(\bar{\tau}_x) = E_p\left[ \sum_{j=1}^{a_j} E_p\left( \frac{\eta_j - \eta_{j-1}}{X_j^{2/3}} \bigg| \bar{X}_j \right) \frac{\bar{X}_j^{2/3}}{x^{2/3}} \right] + \frac{1}{x^{2/3}} E_p[\bar{\tau}_x - \eta_{a_j} \big| R_{a_j}] \] (5.98)

\[ = E_p\left[ \sum_{j=1}^{a_j} E_p\left( \frac{\eta_j}{X_j^{2/3}} \bigg| \bar{X}_1 \right) \frac{\bar{X}_j^{2/3}}{x^{2/3}} \right] + \frac{1}{x^{2/3}} E_p[\bar{\tau}_x - \eta_{a_j} \mid \bar{X}_1 > x - R_{a_j}]. \]

Thus, the proof of (5.96) is a consequence of the 3 inequalities displayed in Claims 5.7, that are also proven below.

Claim 5.7. For every \( \beta > 0 \) there exists a \( C > 0 \) such that

\[ E_p\left[ \frac{\eta_1}{\bar{X}_1^{2/3}} \bigg| \bar{X}_1 = x \right] \leq C, \quad x \in \mathbb{N}, \] (5.99)

\[ E_p\left[ \frac{\bar{\tau}_x}{x^{2/3}} \bigg| \bar{X}_1 > x \right] \leq C, \quad x \in \mathbb{N}, \] (5.100)

\[ E_p\left[ \sum_{j=1}^{a_j} \frac{\bar{X}_j^{2/3}}{x^{2/3}} \right] \leq C, \quad x \in \mathbb{N}. \] (5.101)

Proof. We need to introduce some more notations to prove those three inequalities. We recall (2.13) and we note that, under \( P_p \), the first excursion \( (V_0, \ldots, V_{\eta_1-1}) \) can also be divided into two independent processes, i.e., \( (V_0, \ldots, V_{\eta_1-1}) \) and \( (V_{\eta_1}, \ldots, V_{\eta_{i-1}}) \) with

\[ \tilde{\eta}_1 := \max\{i \geq 1 : V_i = V_{i+1} = \cdots = V_j = 0\}. \] (5.102)

We can therefore rewrite \( X_1 = \tilde{\eta}_1 + Z_1 \) and \( \bar{X}_1 = \tilde{\eta}_1 + \bar{Z}_1 \) with

\[ Z_1 := \eta_1 - \tilde{\eta}_1 + \sum_{i=\tilde{\eta}_1+1}^{\eta_{i-1}} |V_i| \quad \text{and} \quad \bar{Z}_1 := \eta_1 - \tilde{\eta}_1 + \sum_{i=\tilde{\eta}_1+1}^{\eta_{i-1}} V_i - V_{\tilde{\eta}_1}. \] (5.103)

We observe that \( \tilde{\eta}_1 \) and \( \eta_1 \) both follow a geometric law on \( \mathbb{N}_0 \) with parameter \( P_p(U_1 = 0) \) and \( P_p(U_1 \leq 0) \) respectively. Moreover, \( (\tilde{\eta}_1 - \eta_1, Z_1) \) and \( (\eta_1 - \tilde{\eta}_1, \bar{Z}_1) \) have exactly the same law since \( (|V_0|, \ldots, |V_{\eta_1-1}|) \)
and \((V_{\eta_1} - V_{\eta_2}, \ldots, V_{\eta_{n-1}} - V_{\eta_n})\) themselves have the same law which can be seen also as the law of \((V_0, \ldots, V_{n-1})\) under \(P_p(\cdot | V_1 > 0)\).

We will also need the fact that there exists a \(c > 0\) such that
\[
P_p(Y = x) = \frac{c}{x^{1/3}}(1 + o(1)), \quad \text{for} \quad Y \in \{X_1, \tilde{X}_1, Z_1, \tilde{Z}_1\} \quad \text{and} \quad x \in \mathbb{N}.
\] (5.104)

To prove (5.104), we first observe that the local central limit theorem obtained in \((\text{Carmona and Pétrélis, 2016}, \text{Lemma 4.6})\) (which remains true under \(P_p\)) is exactly (5.104) for \(X_1\). From this, we deduce that (5.104) is true with \(Z_1\) as well by using
\[
P_p(Y = x) = \frac{1}{2\pi} \int_0^{2\pi} \phi_Y(t)e^{-ixt} dt, \quad x \in \mathbb{N}, \quad Y \in \{X_1, Z_1\}
\]
(\(\phi_Y\) being the characteristic function of \(Y\)) in combination with the fact that \(X_1 = \tilde{\eta}_1 + Z_1\) and that \(\tilde{\eta}_1\) has a geometric law and is independent of \(Z_1\). As a consequence, (5.104) holds true for \(\tilde{Z}_1\) also since \(Z_1\) and \(\tilde{Z}_1\) have the same law. Finally, the fact that \(\tilde{X}_1 = \tilde{\eta}_1 + \tilde{Z}_1\), that \(\tilde{\eta}_1\) has a geometric law and is independent of \(\tilde{Z}_1\) allows us to conclude that (5.104) is satisfied for \(\tilde{X}_1\) as well.

For the sake of conciseness, we will not display the details of the proof of (5.99). The reason is that the same inequality is proven in \((\text{Carmona and Pétrélis, 2016}, \text{Lemma 4.11})\) with \(\tilde{\eta}_1\) instead of \(\eta_1, \tilde{X}_1\). Then, the equalities \(X_1 = \tilde{\eta}_1 + Z_1\) and \(\tilde{X}_1 = \tilde{\eta}_1 + \tilde{Z}_1\), the equality in law of \((\tilde{\eta}_1 - \eta_1, Z_1)\) and \((\eta_1 - \tilde{\eta}_1, \tilde{Z}_1)\), the fact that \(\tilde{\eta}_1\) and \(\tilde{\eta}_1\) have geometric laws (and therefore light tails) and (5.104) are sufficient deduce (5.99) from \((\text{Carmona and Pétrélis, 2016}, \text{Lemma 4.11})\).

We continue with (5.100). We will assume for simplicity and in this proof only that \(x^{1/3} \notin 2\mathbb{N}\). The case where \(x^{1/3} \notin 2\mathbb{N}\) is taken care of similarly. We use (5.104) to claim that there exists a \(c > 0\) such that \(P_p(\tilde{X}_1 > x) = \frac{c}{x^{1/3}}(1 + o(1))\) therefore, the proof of (5.100) will be complete once we show that there exists a \(C > 0\) such that
\[
E_p(\hat{\tau}_x 1_{\{\tilde{X}_1 > x\}}) \leq C x^{1/3}, \quad x \in \mathbb{N}.
\] (5.105)

We introduce for every \(y \in \mathbb{N}\) the stopping time
\[
\bar{\tau}_y := \inf\{j \geq 1: j + \sum_{i=1}^{j} V_{\eta_i} + t - V_{\eta_i} \geq y\}
\] (5.106)
so that under the event \(\{\tilde{X}_1 > x\}\) we have \(\hat{\tau}_x \leq \tilde{\eta}_1 1_{\{\tilde{\eta}_1 > x\}} + (\tilde{\eta}_1 + \bar{\tau}_{x - \tilde{\eta}_1} 1_{\{\tilde{\eta}_1 < x\}})\) and it allows us to rewrite (5.105) as
\[
E_p(\hat{\tau}_x 1_{\{\tilde{X}_1 > x\}}) \leq E_p(\tilde{\eta}_1) + E_p(\bar{\tau}_{x - \tilde{\eta}_1} 1_{\{\tilde{\eta}_1 < x\}} 1_{\{\tilde{X}_1 > x\}}) := M_{1,x} + M_{2,x}.
\] (5.107)

The first term \(M_{1,x}\) in the r.h.s. of (5.107) is simply bounded above by \(E_p(\tilde{\eta}_1)\) which is finite since \(\tilde{\eta}_1\) has a geometric law. Therefore, it remains to control \(M_{2,x}\) and since \(\tilde{\eta}_1\) is independent of \((V_{\eta_1} - V_{\eta_2}, \ldots, V_{\eta_{n-1}} - V_{\eta_n})\) which (as explained above) has the same law as \((V_i)_{i=0}^{n_1-1}\) under \(P_p(\cdot | V_1 > 0)\) we simply rewrite (recall (2.4))
\[
M_{2,x} = \sum_{k=0}^{n_1-1} P_p(\tilde{\eta}_1 = k) E_p(\bar{\tau}_{x - k} 1_{\{\tilde{X}_1 > x-k\}})
\]
\[
= \sum_{k=0}^{n_1-1} P_p(\tilde{\eta}_1 = k) E_p(\hat{\tau}_{x - k} 1_{\{X_1 > x-k\}} | V_1 > 0).
\] (5.108)

From (5.108), we deduce that the proof of (5.100) will be complete once we show that there exists a \(c > 0\) such that
\[
E_p(\hat{\tau}_x 1_{\{X_1 > x\}} 1_{\{U_1 > 0\}}) \leq cx^{1/3}, \quad x \in \mathbb{N}.
\] (5.109)
To prove (5.109), we write $E_\beta(\xi_x 1_{[\{\xi_x \leq 2x^{1/3}\} \cap \{\xi_x > 0\}]}) \leq \mathcal{F}_{1,x} + \mathcal{F}_{2,x}$ with
\begin{align*}
\mathcal{F}_{1,x} &= E_\beta[\xi_x 1_{[\{\xi_x \leq 2x^{1/3}\} \cap \{\xi_x > 0\}]},
\mathcal{F}_{2,x} &= E_\beta[\xi_x 1_{[\{\xi_x > 2x^{1/3}\} \cap \{\xi_x > 0\}]}].
\end{align*}

The terms $\mathcal{F}_{1,x}$ and $\mathcal{F}_{2,x}$ are taken care of in a similar manner, but proving that there exists a $c > 0$ such that $\mathcal{F}_{2,x} \leq cx^{1/3}$ for every $x \in \mathbb{N}$ is harder than proving the same inequality for $\mathcal{F}_{1,x}$. For this reason and for conciseness, we will only deal with $\mathcal{F}_{2,x}$ here.

We let $L_{x^{1/3}}$ and $\tilde{L}_{x^{1/3}}$ be the last time at which $V$ crosses $x^{1/3}$ before time $\xi_x$ and the first time (after $L_{x^{1/3}}$) at which $V$ crosses $\frac{3}{2}x^{1/3}$, i.e.,
\begin{align*}
L_{x^{1/3}} &= \max\{j \in \{0, \ldots, \xi_x\} : V_j \leq x^{1/3}\},
\tilde{L}_{x^{1/3}} &= \min\{j \geq L_{x^{1/3}} + 1 : V_j \leq \frac{3}{2}x^{1/3}, V_{j+1} > \frac{3}{2}x^{1/3}\}.
\end{align*}

We observe that $\xi_x - L_{x^{1/3}} \leq x^{2/3}$ since after $L_{x^{1/3}}$ the trajectory remains above $x^{1/3}$ up to time $\xi_x$. Therefore, there exists a $C > 0$ such that $E_\beta[\xi_x - L_{x^{1/3}}] \leq Cx^{1/3}$ for every $x \in \mathbb{N}$ and we can safely substitute $L_{x^{1/3}}$ to $\xi_x$ in the definition of $\mathcal{F}_{2,x}$. We define also
\begin{align*}
\mathcal{A}_x := \{(j_1, j_2, y_1, y_2, z_1, z_2) \in (\mathbb{N}_0)^6 : j_1 < j_2 \leq x, \ z_1 \leq z_2 \leq x \text{ and } y_1, y_2 \in \left(x^{1/3}, \frac{3}{2}x^{1/3}\right)\}
\end{align*}
and split the expectation defining $\mathcal{F}_{2,x}$ depending on the values $(j_1, j_2)$ taken by $(T_{x^{1/3}}, \tilde{T}_{x^{1/3}})$ and $(y_1, y_2)$ taken by $(V_{j_1+1}, V_{j_2})$ and also $(z_1, z_2)$ taken by $(K_{j_1}, K_{j_2})$. Thus,
\begin{align*}
\mathcal{F}_{2,x} &= \sum_{(j_1, j_2, y_1, y_2, z_1, z_2) \in \mathcal{A}_x} \sum_{j \geq j_2 \geq x^{1/3}} j_1 \beta(j_1, y_1, z_1) P(j_1, y_1, z_1, j_2, y_2, z_2, j, y_3),
\end{align*}

with
\begin{align*}
W_{j_1, y_1, z_1} &= \{V_j > 0, \forall i \in \{1, \ldots, j_1\}, V_{j_1} \leq x^{1/3}, K_{j_1} = z_1, V_{j_1+1} = y_1\},
I_{j_1, y_1, z_1, j_2} &= \{x^{1/3} < V_{j_1} \leq \frac{3}{2}x^{1/3}, \forall i \in \{j_1 + 1, \ldots, j_2\}, V_{j_2} = y_2, K_{j_2} = z_2\},
J_{j, y_1, z_1, j_2} &= \{V_{j_2+1} = y_3, V_j > x^{1/3}, \forall i \in \{j_2 + 1, \ldots, \xi_x\}, V_{\xi_x} > \frac{3}{2}x^{1/3}\}.
\end{align*}

Provided we change the equality into an inequality, we can safely restrict the event $J_{j_2}$ in the r.h.s. of (5.112) to $\{V_{j_2+1} = y_3\}$. Then, we apply Markov property at time $j_1 + 1$ and $j_2$ and we obtain
\begin{align*}
\mathcal{F}_{2,x} \leq \sum_{(j_1, j_2, y_1, y_2, z_1, z_2) \in \mathcal{A}_x} \sum_{j \geq j_2 \geq x^{1/3}} j_1 \beta(j_1, y_1, z_1) P(j_1, y_1, z_1, j_2, y_2, z_2) P(j_2, y_2)(V_1 = y_3),
\end{align*}

with
\begin{align*}
\tilde{I}_{j_1, j_2, y_1, y_2, z_1, z_2} &= \{x^{1/3} < V_{j_1} \leq \frac{3}{2}x^{1/3}, \forall i \leq j_2 - j_1 - 1, V_{j_2-1} = y_2, K_{j_2-1} = z_2 - z_1 - 1 - y_1\}.
\end{align*}

At this stage, we map every trajectory $(V_i)_{i=1}^{j_2+1}$ taken account in $W_{j_1, y_1, z_1}$ onto an associated path which equals $V$ up to time $j_1$, then touches $x^{1/3}$ at time $j_1 + 1$ and is equal to $2x^{1/3} - y_1$ at time $j_2 + 2$ (i.e. we reflect $V_{j_2+1}$ with respect to $x^{1/3}$ to obtain the position of the image of $V$ at time $j_1 + 2$). Thus, we obtain a new set
\begin{align*}
\tilde{W}_{j_1, y_1, z_1} &= \{V_j > 0, \forall i \in \{1, \ldots, j_1\}, V_{j_1} \leq x^{1/3}, K_{j_1} = z_1, V_{j_1+1} = x^{1/3}, V_{j_2+2} = 2x^{1/3} - y_1\},
\end{align*}
so that $P_\beta(W_{j_1, y_1, z_1}) = c_\beta P_\beta(\tilde{W}_{j_1, y_1, z_1})$. We also reflect every piece of trajectory in $\tilde{I}_{j_1, j_2, y_1, y_2, z_1, z_2}$ with respect to $x^{1/3}$ and we denote by $\tilde{I}_{j_1, j_2, y_1, y_2, z_1, z_2}$ the set containing the resulting paths, thus
\begin{align*}
P_\beta(j_1, y_1, z_1, j_2, y_2, z_2) = P_\beta(2x^{1/3} - y_1, j_1, j_2, y_1, y_2, z_1, z_2).
\end{align*}
and
\[ \tilde{t}_{j_i,j_{i+1}} \subset \{ \frac{1}{2} x^{1/3} \leq V_i < x^{1/3}, \forall i \leq j_2 - j_1 - 1, V_{j_2-j_1-1} = 2x^{1/3} - y_2, K_{j_2-j_1-1} \leq z_2 - z_1 \}. \]

With the help of (5.116–5.117) and by summing over $y_3$ in (5.114) we obtain
\[
\mathcal{P}_{2,x} \leq c \beta \sum_{(j_1,j_2,j_3) \in E} j_1 \mathbb{P}(W_{j_1,j_3}) \mathbb{P}(U_1 > \frac{3}{2} x^{1/3} - y_2). 
\]  
(5.118)

Since $U_1$ has a geometric law, there exists a $c > 0$ such that $\mathbb{P}(U_1 > y) = c \mathbb{P}(U_1 = y)$ for every $y \geq 0$. Therefore, there exists a constant $c > 0$ such that we can substitute $\mathbb{P}(U_1 = y_2 - \frac{3}{2} x^{1/3})$ to $\mathbb{P}(U_1 > \frac{3}{2} x^{1/3} - y_2)$ in the r.h.s. in (5.118) and write
\[
\mathcal{P}_{2,x} \leq c \sum_{(j_1,j_2,j_3) \in E} j_1 \mathbb{P}(W_{j_1,j_3}) \mathbb{P}(U_1 > \frac{3}{2} x^{1/3} - y_2) \cap (V_{j_2-j_1} = \frac{1}{2} x^{1/3})). 
\]  
(5.119)

At this stage, we let $C_{j_1,j_2,j_3}$ be the set of paths obtained by concatenating trajectories in $W_{j_1,j_3}$ and in $\tilde{t}_{j_1,j_2,j_3} \cap (V_{j_2-j_1} = \frac{1}{2} x^{1/3})$, i.e.,
\[
C_{j_1,j_2,j_3} \subset \{ V_i > 0, \forall i \in \{1, \ldots, j_1\}, V_i \leq x^{1/3}, K_i = z_1, V_{j_1} = x^{1/3}, V_{j_1+1} = 2x^{1/3} - y_1 \}
\[ \frac{1}{2} x^{1/3} \leq V_i < x^{1/3}, \forall i \in \{2 + j_1, \ldots, j_2\}, V_{j_2+1} = 2x^{1/3} - y_2, \]
\[ K_{j_2+j_1} \leq z_2 + x^{1/3}, V_{j_2+j_1} = \frac{1}{2} x^{1/3} \}, 
\]  
(5.120)

and it is fundamental to note the $C_{j_1,j_2,j_3}$ with $(j_1,j_2,j_1,j_2,1) \in E_x$ are disjoint. The final step of this proof consists in attaching at the end of every path in a given $C_{j_1,j_2,j_3}$ another path which will reach the lower half-plane in such a way that the area swept by the whole excursion belongs to $[\frac{1}{2}, 2x]$. We continue this computation by noticing again that by Doob’s Theorem
\[
\lim_{x \to \infty} \mathbb{P}(V_i < x^{1/3}, \forall i \leq \tau, K_i \in [\frac{1}{2}, x]) = \mathbb{P}_{\beta=\frac{1}{2}}(B_0 < 1, \forall s \in [0, \tau_0], A_s \in [\frac{1}{2}, 1]) > 0, 
\]  
(5.121)

and therefore we can assert that there exists a $c > 0$ such that for every $x \in \mathbb{N}$, it holds that
\[
\mathcal{P}_{2,x} \leq c \sum_{(j_1,j_2,j_3) \in E} j_1 \mathbb{P}(C_{j_1,j_2,j_3}) \mathbb{P}(V_i < x^{1/3}, \forall i \leq \tau, K_i \in [\frac{1}{2}, x]) 
\]  
(5.122)

and
\[
\mathcal{P}_{2,x} \leq c \sum_{(j_1,j_2,j_3) \in E} j_1 \mathbb{P}(C_{j_1,j_2,j_3}) \mathbb{P}(V_i < x^{1/3}, \forall i \leq \tau, K_i \in [\frac{1}{2}, x]) \leq c \mathbb{E}_p \left[ \max\{1_{\{V_i \geq x^{1/3}\}} \right] 
\]  
with $\max\{1_{\{V_i \geq x^{1/3}\}} := \max(1 \leq i \leq N : V_i \geq x^{1/3})$. This provides us the expected upper bound on $\mathcal{P}_{2,x}$ and the proof of (5.100) is complete.

We conclude this section with the proof of (5.101). To begin with we recall the definition of $\alpha_x$ in (5.93) and we note that by definition of $\alpha_x$
\[
\mathbb{E}_p \left[ \sum_{j=1}^{a_x} \bar{X}_{2j} \right] = \mathbb{E}_p \left[ \sum_{j=1}^{a_x} \bar{X}_{2j} 1_{\{\bar{X}_{2j} \leq x\}} \right] \leq \mathbb{E}_p \left[ \sum_{j=1}^{a_x} \bar{X}_{2j} 1_{\{\bar{X}_{2j} \leq x\}} \right]. 
\]  
(5.123)

Since $1 + \alpha_x$ is a bounded stopping time with respect to the filtration $(\sigma(\bar{X}_1, \ldots, \bar{X}_n))_{n \geq 1}$ and since $(\sum_{j=1}^{a_x} \bar{X}_1 1_{\{\bar{X}_{2j} \leq x\}} - n \mathbb{E}_p(\bar{X}_1 1_{\{\bar{X}_{2j} \leq x\}}))_{n \geq 1}$ is a martingale, we can rewrite (5.123) as
\[
\mathbb{E}_p \left[ \sum_{j=1}^{a_x} \bar{X}_{2j} 1_{\{\bar{X}_{2j} \leq x\}} \right] \leq \mathbb{E}_p \left[ \sum_{j=1}^{a_x} \bar{X}_{2j} 1_{\{\bar{X}_{2j} \leq x\}} \right]. 
\]  
(5.124)

A straightforward computation with the help of (5.104) guarantees that there exists a $c > 0$ such that $\mathbb{E}_p(\bar{X}_{2j} 1_{\{\bar{X}_{2j} \leq x\}}) \leq cx^{1/3}$ for every $x \in \mathbb{N}$. Then, it remains to use (5.104) for $\bar{X}_1$ to conclude that exists a $c > 0$ such that $\mathbb{E}(\alpha_x) \leq cx^{1/3}$ for every $x \in \mathbb{N}$. 
\]
5.3. Proof of Proposition 2.5

We shall prove first the second limit. Thanks to the independence of $D$ and $B$ the process

$$ N(s) := D(s) - 	ilde{B}^k(s) = \int_0^s 1_{(u \notin \Gamma_k)} \, dD_u $$

is an $L^2$ martingale, and by Doob maximal inequality

$$ \mathbb{P} \left( \sup_{x \in [0, 1]} |N(s)| \geq \varepsilon \right) \leq \frac{C}{\varepsilon^2} \mathbb{E} \left[ N(1)^2 \right] $$

$$ = \frac{C}{\varepsilon^2} \int_0^\infty \mathbb{P}(u \notin \Gamma_k, u < a) \, du $$

Since the excursion process of Brownian motion is sigma-finite, we have for all $u : \mathbb{P}(u \notin \Gamma_1) \to 0$. Therefore we can conclude by dominated convergence if we can prove that $\mathbb{E}[a_1] < +\infty$. This is indeed true since $Y_t = |B|_{\delta(t)}$ conditioned by $Y(1) = 0$ is distributed as $\left( \frac{2}{\delta} \rho \right)^{2/3}$ with $(\rho_t)_{0 \leq t \leq 1}$ a Bessel bridge of dimension $\delta = 4/3$. Therefore,

$$ \mathbb{E}[a_1] = \mathbb{E} \left[ \int_0^1 \frac{ds}{Y_s} \right] = C \int_0^1 \mathbb{E}[\rho_t^{-2/3}] \, dt $$

Since $\rho_t$ has density

$$ C_\delta(t(1-t))^{-\delta/2} \exp(-\frac{x^2}{2} + \frac{1}{1-t}) $$

we see that, by symmetry,

$$ \mathbb{E}[a_1] \leq C \int_0^1 t^{-\delta/2} \left( \int_{x=0}^\infty x^{-2/3} e^{-\frac{x^2}{2}} \right) \, dt < +\infty. $$

We now prove the first limit of Proposition 2.5. Given $\eta > 0$, since $\mathbb{E}[a_1] < +\infty$ there exists $A = A(\eta)$ such that

$$ \mathbb{P}(a_1 \geq A) \leq \frac{1}{2} \eta. $$

By Kolmogorov’s Continuity Criterion given $0 < \gamma < \frac{1}{2}$ the random variable

$$ K := \sup_{0 \leq s < t \leq A} \frac{|B_t - B_s|}{|t - s|^\gamma} $$

has a small moment: there exists $\delta > 0$ such that $\mathbb{E}[K] < +\infty$.

If $s \leq 1$ then, on $\{a_1 < A\}$,

$$ |\tilde{B}(s) - \tilde{B}^k(s)| = |B(a_s)| 1_{(a_s, \notin \Gamma_k)} $$

$$ = |B(a_s - B(a_s, s)| 1_{(a_s, \notin \Gamma_k)} $$

$$ \leq |a_s - g(a_s)| \, K 1_{(a_s, \notin \Gamma_k)} $$

$$ \leq \frac{1}{k^\gamma} K $$

Therefore

$$ \mathbb{P} \left( \sup_{x \in [0, 1]} |\tilde{B}(s) - \tilde{B}^k(s)| \geq \varepsilon \right) \leq \mathbb{P}(a_1 \geq A) + \mathbb{P} \left( \sup_{x \in [0, 1]} |\tilde{B}(s) - \tilde{B}^k(s)| \geq \varepsilon; a_1 < A \right) $$

$$ \leq \frac{1}{2} \eta + \mathbb{P}(K \geq k^{\gamma} \varepsilon) \leq \frac{1}{2} \eta + \frac{1}{(k^{\gamma} \varepsilon)^\delta} \mathbb{E}[K^\delta] < \eta, $$

for $k$ large enough.
References

BILLINGSLEY, P. (2008). Convergence of Probability Measures. John Wiley and Sons, Inc., New York. MR1700749 (2000e:60008)

BRAK, R., GUTTMANN, A. J. and WHITTINGTON, S. G. (1992). A collapse transition in a directed walk model. J. Phys. A 25 2437–2446. MR1164357 (93a:82025)

CRAVENNA, F and DUESCHEL, J.-D. (2009). Scaling limits of (1 + 1)-dimensional pinning models with Laplacian interaction. Ann. Probab. 37 903–945. MR2537545 (2010j:60237)

CRAVENNA, F, SUN, R. and ZYGOURAS, N. The continuum disordered pinning model. http://arxiv.org/abs/1406.5088 (math.PR), submitted.

CARMONA, P., NGUYEN, G. B. and PETRÉLIS, N. (2016). Interacting partially directed self avoiding walk. From phase transition to the geometry of the collapsed phase. Ann. Probab. 44 3234–3290.

CARMONA, P. and PETRÉLIS, N. (2017). Limit theorems for random walk excursion conditioned to have a typical area. https://arxiv.org/abs/1709.06448.

DENISOV, D., KOLB, M. and WACHTEL, V. (2015). Local asymptotics for the area of random walk excursions. J. Lond. Math. Soc. (2) 91 495–513. MR3355112

DEUSCHEL, J. D., GIACOMIN, G. and ZAMBOTTI, L. (2005). Scaling limits of equilibrium wetting models in (1 + 1)-dimension. Probab. Theory Related Fields 132. MR2198199 (2007f:60080)

DURRETT, R. (2010). Probability: theory and examples, fourth ed. Cambridge Series in Statistical and Probabilistic Mathematics. Cambridge University Press, Cambridge. MR2722836 (2011e:60001)

NGUYEN, G. B. and PETRÉLIS, N. (2013). A variational formula for the free energy of the partially directed polymer collapse. J. Stat. Phys. 151 1099–1120. MR3063498

SOHIER, J. (2013). The scaling limits of a heavy tailed Markov renewal process. Ann. Inst. H. Poincaré Probab. Statist. 49 483–505.

ZWANZIG, R. and LAURITZEN, J. I. (1968). Exact calculation of the partition function for a model of two dimensional polymer crystallization by chain folding. J. Chem. Phys. 48 3351.