Growth and fabrication of InAs/GaSb type II superlattice mid-wavelength infrared photodetectors

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Abstract
We report our recent work on the growth and fabrication of InAs/GaSb type II superlattice photodiode detectors. The superlattice consists of 9 monolayer InAs/12 monolayer GaSb in each period. Lattice mismatch between the GaSb substrate and the superlattice is $1.5 \times 10^{-4}$. The full width at half maximum of the first-order satellite peak from X-ray diffraction is 28 arc sec. The P-I-N photodiodes in which the absorption regions (I regions) have 600 periods of superlattice show a 50% cutoff wavelength of 4.3 $\mu$m. The current responsivity was measured at 0.48 A/W from blackbody radiation. The peak detectivity of $1.75 \times 10^{11}$ cmHz$^{1/2}$/W and the quantum efficiency of 41% at 3.6 $\mu$m were obtained.

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Keywords: InAs/GaSb, type II superlattice, photodiodes, infrared

Introduction
HgCdTe [MCT] photodetectors which offer excellent quantum efficiency are the dominating infrared technology, and very large MCT sensor arrays are available. The drawbacks of MCT come from its technology difficulty. MCT has weak mechanical strength due to the weak ionic bonds and low uniformity due to the high Hg vapor pressure. Common substrates for MCT epitaxial growth are lattice-matched CdZnTe or readily available Si or Ge capped with a few-micron-thick buffer layers, yet no substrates are known to date which can satisfy all necessities for being low cost, lattice-matched, and chemically, mechanically, and optically well suited. Thus, the fabrication of large-format MCT arrays with homogeneous performances becomes more and more challenging, especially with an increasing cutoff wavelength. Alternative approaches based on quantum, mechanically tailored semiconductor heterostructures have been developed including quantum well infrared photodetectors, quantum dot infrared photodetectors, and InAs/GaSb type II, strained-layer superlattice [SLS] photodetectors.

The idea of the InAs/GaSb superlattice [SL] was firstly introduced by Sai-Halasz et al. in 1977 [1]. It was proposed to be used for infrared detection by Smith and Mailhoit in 1987 [2]. Since then, the InAs/GaSb SL has received more and more attentions for infrared detection due to its unique advantages over other infrared materials. The effective bandgap of the SL material can be tailored over a wide range ($3 \mu m \leq \lambda \leq 30 \mu m$) by varying the thickness of InAs and GaSb, the two 'mid-gap' constituent materials. Auger recombination is suppressed due to large splitting between the heavy-hole and light-hole subbands induced by the strains in the SL. Tunneling currents in SL detectors are reduced due to the large electron effective mass. The band structures of the SL material can be engineered to enhance carrier lifetimes and reduce noises. InAs/GaSb SL detectors are actually the only candidate that has theoretically predicted performances better than MCT [3-5]. Moreover, the type II SLs have the advantages of excellent uniformity and low cost due to the mature III-V material technology. The first high-performance InAs/GaSb photodetector was demonstrated in 1996 [6]. The first SL focal plane array [FPA] detector was reported in
2003 [7]. Some leading laboratories have recently demonstrated megapixel FPA detectors [8,9].

Despite the rapid progress, InAs/GaSb type II SL photodetectors are still at their infancy time. Their performances have far not reached their theoretical prediction yet and are even inferior to those of MCT. One important issue which limited the performances of SLS detectors is the SL material quality [10]. Yang et al. reported that the carrier lifetimes of SLS are limited by the Shockley-Read-Hall recombination when the carrier concentration is lower than 10^{17} cm^{-3} [11]. Therefore, it is desired to grow high-quality SL materials with high crystal perfection for device applications. Rodriguez et al. reported their mid-wavelength infrared SL materials with an X-ray diffraction full width at half maximum [FWHM] of 26 arc sec [12]. Khoshakhlagh et al. [13] reported a FWHM of about 32 arc sec for a SL material with an 8-μm cutoff wavelength. In SL material growth, the control of interface type and quality dramatically affects the overall material quality [13]. As InAs and GaSb have no common atoms and both arsenic and antimony’s sticking coefficients are less than 1, two types of interfaces may be formed, the GaAs-like and the InSb-like, according to the growth conditions. Therefore, in InAs/GaSb SL growth, the interface controls are particularly complicated and important.

We report in this study the material growth and device fabrication of InAs/GaSb SL mid-infrared photodiodes. In particular, we designed a new shutter sequence in the growth process to improve the interface quality. This is our first effort to bring up the InAs/GaSb material technology for infrared detection in our laboratory.

**Material growth and device fabrication**

The InAs/GaSb SL is grown by molecular beam epitaxy [MBE] on n-type doped (100) GaSb substrates. The growth temperature was set at 450°C and a V/III ration of 5:1. There are two important issues uniquely associated with InAs/GaSb SL growth: (1) InAs has 0.75% smaller lattice constant than GaSb, and proper interface layers, in typical InSb, have to be inserted between the InAs and GaSb layers for strain balance and (2) interface control is extremely important to obtain high-quality epitaxial materials since there are no common atoms between InAs and GaSb. Shutter sequences were carefully designed to favor an InSb-like interface layer with a desired thickness. InSb-like interfaces were realized through proper shutter sequences. An important issue is to suppress the arsenic flux at the interface growth since arsenic has a high background pressure in a MBE chamber. At GaSb-to-InAs interfaces, we first closed the gallium cell shutter, left the antimony cell shutter open for 2 s, and then opened the indium cell for 0.8 s before opening the arsenic shutter for the InAs growth. At the InAs-to-GaSb interface, we first closed the arsenic shutter and left the indium shutter open for 1 s to have the growing surface covered by one layer of indium atoms, then closed the indium shutter, and opened the antimony shutter to switch to the GaSb growth. With the above shutter sequence, we greatly reduced the arsenic background pressure during the InSb and GaSb growth. The SL crystal qualities, such as the period and the lattice-mismatch from the GaSb substrates, were characterized by high-resolution X-ray diffraction [HRXRD]. The surface quality was measured by atomic force microscopy [AFM].

The photodetector has a P-I-N structure. It consists of 320-nm-thick p-type SLs with a 10^{17}.cm^{-3} Be doped in the GaSb layer, 3.2-μm-thick undoped i-type SLs, and 320-nm-thick n-type SLs with a 1.5 × 10^{17} cm^{-3} Si doped in the InAs layer. All the SLs consist of 9 mono-layer [ML] InAs and 12 ML GaSb. A 1-μm-thick GaSb layer with Be doping to an order of 5 × 10^{17} cm^{-3} and a 30-nm-thick InAs layer with Si doping to an order of 5 × 10^{17} cm^{-3} were below and above the P-I-N structure, respectively, as a contact layer. The single-element detectors have architecture as shown in Figure 1. The detectors are designed to receive the irradiance from the front side in order to avoid the strong GaSb substrate absorption [10].

Contact photolithography was employed for device fabrication. The mesas were wet etched with a mixed solution of citric acid, phosphoric acid, and hydrogen peroxide (10:1:1). The sidewall of the mesa is then passivated by sputtering a 300-nm-thick SiO2 layer. Figure 2 shows a scanning electron micrograph [SEM] of an
etched and passivated detector mesa. Contact windows were then opened by inductively coupled plasma reactive-ion etching. A composite contact layer of 20-nm Ti, 30-nm Pt, and 20-nm Au was deposited by electron-beam evaporation on both the p-type GaSb buffer layer and the n-type top InAs layer. It was followed by thermal vapor evaporation of a thick Au layer for wire bonding. Finally, single-element devices of different photosensitive areas varying from $100 \mu m \times 100 \mu m$ to $500 \mu m \times 500 \mu m$ were fabricated.

The detectors were mounted onto homemade cold fingers in a Dewar, which were cooled to 77 K with liquid nitrogen. Spectral responsivity was measured using Fourier transform infrared (FTIR) spectroscopy. The current-voltage and dynamic resistance (DR)-voltage curves were swept by a Keithley 236 source-measure unit (Keithley Instruments, Inc., Shanghai, China) using a self-coded LabVIEW program. For photoresponse measurements, the blackbody temperature was set at 800 K, and the chopper, at 800 Hz. The signals were picked up by a pre-amplifier and a lock-in amplifier.

**Results and discussion**

Figure 3 shows a high-resolution $\omega$-$2\Theta$ scanning curve of a 9 ML InAs/12 ML GaSb SL. The layer numbers of InAs and GaSb in each period were determined to achieve a 4.5-μm cutoff wavelength using a k-p model under the envelope function approximation. The SL consists of 100 periods. Clear and sharp satellite peaks up to the fourth order are observed. Higher-order peaks could be observed if we had extended the scanning angle range. The measured lattice mismatch between the SL and the GaSb substrate is $\Delta a/a = 1.5 \times 10^{-4}$. The FWHM of the first-order peak is 28 arc sec. The thickness of each period is 65.3 Å according to the simulation. It has only a 2% difference comparing with the design thickness (63.8 Å).

The surface morphology was studied with a Digital Instruments Nanoscope atomic force microscope (Santa Barbara, CA, USA) as shown in Figure 4. Atomic steps are clearly observed. We have achieved a root-mean-square roughness of 1.5 Å over an area of 2 μm $\times$ 2 μm.

Figure 5 shows the measured dark current curve and the corresponding DR curve of a 200 μm $\times$ 200 μm
photodiode. For this photodiode, a resistance-area product at zero bias $[R_o A]$ of 147 $\Omega$ cm$^2$ was measured. A current responsivity of 0.48 A/W was measured for the same device, and it has a blackbody detectivity of $4.54 \times 10^{10}$ cmHz$^{1/2}$/W. The InAs/GaSb SL detectors have a 5% cutoff wavelength of 4.3 $\mu$m at a longer wavelength side and 2.0 $\mu$m at a shorter wavelength side according to the FTIR spectroscopy. Combining the response spectrum and the blackbody current responsivity, the absolute current responsivity spectrum and quantum efficiency can be calculated. The result is shown in Figure 6. The numbers labeled in the figure represent the quantum efficiency. The $G$ factor was calculated to be 0.26, and the peak detectivity reaches to $1.75 \times 10^{11}$ cmHz$^{1/2}$/W.

Conclusions

In summary, we grew and fabricated InAs/GaSb type-II SL materials and devices by MBE and wet chemical etching. Lattice mismatch between the substrate and the SL is $1.5 \times 10^{-4}$. The FWHM of the first-order satellite peak from XRD is 28 arc sec. The mid-wavelength infrared photodiodes have a cutoff wavelength of 4.3 $\mu$m. A current responsivity of 0.48 A/W and a peak detectivity of $1.75 \times 10^{11}$ cmHz$^{1/2}$/W were measured. The quantum efficiency of the device at 3.6 $\mu$m is 41%.

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Authors’ contributions

JW contributed the main ideas of SL structure design and supervised the MBE growth. QX carried out the MBE growth. YZ carried out the SL X-ray measurements. JJ carried out the device processing. CL supervised the device processing and carried out the device measurements. LH initiated and supervised the SL infrared detector program. All authors read and approved the final manuscript.

Competing interests

The authors declare that they have no competing interests.

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