Stellar tidal disruption candidates found by cross-correlating the ROSAT Bright Source Catalogue and XMM-Newton observations

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ABSTRACT

We performed a systematic search for stellar tidal disruption events (TDE) by looking for X-ray sources that were detected during the ROSAT All Sky Survey and faded by more than an order of magnitude over the next two decades according to XMM-Newton serendipitous observations. Besides a number of highly variable persistent X-ray sources (like active galactic nuclei and cataclysmic variables), we found three sources that are broadly consistent with the TDE scenario: 1RXS J114727.1+494302, 1RX J130547.2+641252, and 1RXS J235424.5-102053. A TDE association is also acceptable for the fourth source, 1RXS J112312.7+012858, but an AGN origin cannot be ruled out either. This statistics implies a TDE rate of \(~3 \times 10^{-5}\) yr\(^{-1}\) per galaxy in the Universe within \(z \sim 0.18\), which is broadly consistent with the estimates of the TDE rate in the more local Universe obtained previously.

Key words: accretion, accretion discs – black hole physics – methods: observational – galaxies: nuclei.

1 INTRODUCTION

Tidal disruption of a star by a supermassive black hole (SMBH) is expected to be a rather rare event, taking place every \(10^3 \sim 10^6\) years in a typical galaxy in the local Universe (Wang & Merritt 2004). According to the canonical picture, some fraction (~10%), (Ayal, Livio, & Piran 2004) of the disrupted star’s material should be captured and then accreted by the SMBH on a time-scale of about a year (Gurzadian & Ozerov 1981; Rees 1988; Evans & Kochanek 1989; Phinney 1989; Ulmer 1999). This results in the formation of an accretion disc emitting quasi-thermally with the peak luminosity of \(10^{35} \sim 10^{39}\) erg s\(^{-1}\), mainly at extreme ultraviolet (EUV)/soft X-ray wavelengths (Strüder & Oudata 2003). Thus, a previously EUV/X-ray faint galaxy can become active for some period of time, whilst the stellar debris are being accreted onto the SMBH, and then return to its initial quiescent state. Such tidal disruption events (TDE) can be searched for using repeated wide-area surveys spaced in time by at least a year (Sembay & West 1993; Komossa 2002; Gezari et al. 2009; van Velzen et al. 2011b; Stone, Sari, & Loeb 2013; Khabibullin, Sazonov, & Sunyaev 2014).

Only about a dozen TDE candidates have been identified so far (see Komossa 2012; Gezari 2012 for recent reviews), with the major contribution coming from X-ray observations (Komossa 2002; Donley et al. 2002; Esquej et al. 2008; Cappelluti et al. 2009; Maksym, Ulmer, & Eracleous 2010; Lin et al. 2011; Saxton et al. 2012; Maksym et al. 2013). The ROSAT All-Sky Survey (RASS, Voges et al. 1999) played the key role in the first X-ray identifications and provided a reference point for comparison with consequent ROSAT pointed observations (Komossa 2002; Donley et al. 2002) and later observations by Chandra and XMM-Newton (e.g., Esquej et al. 2008). Although up to several thousand TDEs had been expected to be detected during the RASS (Sembay & West 1993), only five were actually identified due to the lack of deep archival and timely follow-up X-ray observations needed for probing the pre- and/or post-outburst state (Komossa 2002; Donley et al. 2002). Therefore, plenty of TDEs will likely remain hidden in the RASS data at least until the forthcoming X-ray all-sky survey by the eROSITA telescope of the SRG observatory (eRASS, Khabibullin, Sazonov, & Sunyaev 2014).

Nevertheless, currently operating X-ray observatories have already covered a significant fraction of the sky with sensitivity at least ten times better than that of the RASS (~3×10\(^{-13}\) erg s\(^{-1}\) cm\(^{-2}\)) in the 0.5–2 keV energy band, (Brandt & Hasinger 2005), thus providing an opportunity to search for TDEs that occurred during the RASS epoch (in 1990–1991) and have decayed since then. In the present study, we check all sources from the RASS Bright Source Catalogue (RASS-BSC, Voges et al. 1999) for a large (more than a factor of ten) flux decrease in \(XMM-Newton\) observations carried out between February 2000 and December 2012, which have covered ~2% of the sky with a characteristic detection limit of \(~10^{-14}\) erg s\(^{-1}\) cm\(^{-2}\) in the 0.5–2 keV energy band.
Similarly to the present work, [Donley et al. 2002] looked for X-ray sources that occurred during the RASS. The difference is that we use the much deeper XMM-Newton serendipitous survey instead of ROSAT pointed observations for comparison with the RASS data. This allows us to search for fainter TDEs, down to the flux limit of RASS-BSC, in the fields observed by XMM-Newton. As a result, although the total area covered by XMM-Newton observations is much smaller than that covered by ROSAT pointed observations, we can probe the TDE rate out to a larger distance (\(z \approx 0.2\)) compared to Donley et al. [2002], while covering a similar volume of the Universe. Comparison of RASS and XMM-Newton data has previously been done by [Esquej et al. 2008]. However, their study was based on the XMM-Newton Slew Survey, whose sensitivity and sky coverage are similar to the survey constructed from ROSAT pointed observations, and it was aimed at flares that occurred during the XMM-Newton epoch and were undetected by the RASS. Due to the relatively small volume of the Universe probed by that study, it discovered only two TDE candidates, which nonetheless proved to be of high value thanks to their timely follow-up observations enabled by the accurate XMM-Newton localisations [Esquej et al. 2008].

In summary, with this study we aim to significantly increase the available sample of TDE candidates and to begin to find TDEs outside the local (\(z < 0.1\)) Universe. In the latter regard, the present study is similar to the much larger forthcoming eROSITA survey [Khabibullin, Sazonov, & Sunyaev 2014]. An obvious drawback of the present study is the impossibility of follow-up observations of TDEs detected during the RASS due to their poor localisation and disappearance from the sky by the present day.

The outline of the paper is as follows. In Section 2 we review the databases and methods used in our analysis. A summary of the identified TDE candidates is presented in Section 3. In Section 4 we discuss the implications of this study for the TDE rate in the Universe. Some conclusions are provided in Section 5.

2 METHOD

Our study follows a number of previous searches for X-ray sources demonstrating a large amplitude flux drop (> 10) and a soft, approximately \(\sim 0.05\) keV blackbody spectrum expected of TDEs [Komossa 2002; Donley et al. 2002; Esquej et al. 2008; Cappelluti et al. 2009; Maksym, Ulmer, & Eracleous 2010; Lin et al. 2011; Saxton et al. 2013; Maksym et al. 2013]. Namely, we looked for sources that were bright in the RASS epoch and faded away over the next \(\sim 10–20\) years, as revealed by serendipitous observations with XMM-Newton, since we do not expect any significant TDE-related emission at this very late phase.

2.1 ROSAT Bright Source Catalogue

The ROSAT All-Sky Survey Bright Source Catalogue (RASS-BSC) was constructed from the all-sky survey performed during the first six months of the ROSAT mission in 1990/91 [Voges et al. 1999]. It contains 18,811 sources with a limiting count rate of 0.05 cts/s in the 0.1–2.4 keV energy band, which corresponds to at least 15 photons from a source. The typical positional accuracy is 30 arcsec and at a brightness limit of 0.1 cts/s (8,547 sources), the catalogue provides a sky coverage of 92%. The catalogue also provides the source extent, i.e. the amount by which the image of a source exceeds the point spread function, which could be used to distinguish extended and point sources.

The total energy band is divided in four channels: A(\(-0.1–0.4\) keV), B(\(-0.5–2.0\) keV), C(\(-0.5–0.9\) keV) and D(\(-0.9–2.0\) keV), and spectral characteristics are presented in the form of hardness ratios H1=(CR(B)- CR(A))/(CR(B)+CR(A)) and H2=(CR(D)- CR(C))/(CR(D)+CR(C)), where CR(X) is the count rate in a given channel X.

There are two standard methods for RASS-BSC [Voges et al. 1999] to convert measured count rates to the unabsorbed 0.1–2.4 keV flux, which differ in the assumptions about the source spectral shape. One assumes an absorbed power law with a fixed photon index \(\Gamma = 2.3\), which is the typical value derived from ROSAT observations of extragalactic sources, and the absorbing column density fixed at the Galactic value along the given line of sight (Flux 1). The other approach (Flux 2) is based on the empirical relation for count rates and fluxes originally suited for stars [Voges et al. 1999]. For these methods, the accepted detection limit of 0.05 cts/s translates to the following 0.1–2.4 keV flux limits for point-sources: Flux1 \(\sim 7 \times 10^{-13}\) erg s\(^{-1}\) cm\(^{-2}\) (assuming \(N_H = 10^{20}\) cm\(^{-2}\)) and Flux2 \(\sim 2 \times 10^{-13}\) erg s\(^{-1}\) cm\(^{-2}\) (for HR1=1–1, see the conversion formula in Section 3.3.5 of [Voges et al. 1999]).

We make use of the FITS file with the catalogue available at the anonymous ftp-server ftp.xray.mpe.mpg.de\(^4\). The analysis was performed using standard general purpose tasks for the manipulation of FITS data contained in the the sub-package FUTILS of the HEASARC’s FTOOLS package\(^3\) [Blackburn et al. 1995].

2.2 XMM-Newton Serendipitous Source Catalogue

3XMM-DR4\(^4\) is the third generation catalogue of X-ray sources serendipitously detected by the European Space Agency’s XMM-Newton observatory. It follows the previous 2XMM catalogue [Watson et al. 2009, 2013].

The current version contains source detections drawn from 7427 XMM-Newton EPIC observations made between 2000 February 3 and 2012 December 8, covering a total of 794 deg\(^2\) of the sky, accounting for overlaps. The position for a typical source can be determined with \(\sim 2\) arcsec accuracy [Watson et al. 2004]. The median flux in the soft (0.2–2 keV) energy band is \(\sim 6 \times 10^{-15}\) erg s\(^{-1}\) cm\(^{-2}\), which is \(\sim 100\) times lower than the detection threshold accepted for RASS-BSC (see above).

In this study we are interested not only in X-ray sources detected by XMM-Newton but also in X-ray flux limits provided by XMM-Newton observations. Hence, we use the XMM flux upper limit server (FLIX\(^3\)), which makes it possible, by scanning the XMM-Newton images corresponding to the 3XMM-DR4 catalogue, to estimate an upper limit on the X-ray flux from a given position as well as to measure the flux within a circle centred on that position.

2.3 Revealing decay of RASS-BSC sources with XMM-Newton observations

Our goal is to obtain a sample of TDE candidates by cross-correlating the RASS-BSC source catalogue with the XMM-Newton data. As usual in such studies, we should take both reliability and completeness of the resulting sample into account. We

\(^{1}\) Or directly at [ftp://ftp.xray.mpe.mpg.de/rosat/catalogues/rass-bsc/]

\(^{2}\) http://heasarc.gsfc.nasa.gov/ Photolab/

\(^{3}\) http://xmmssc-www.star.le.ac.uk/Catalogue/xcat_public_3XMM-DR4.html

\(^{4}\) http://www.ledas.ac.uk/flux/flux3
prefer to use a more reliable selection algorithm, since it is virtually impossible to perform follow-up observations to confirm or reject a TDE association for RASS sources.

As a first step, we leave only sources outside the Galactic plane, by requiring $|b| > 30^\circ$. There are two reasons for applying this constraint. First, the Galactic extinction at $|b| < 30^\circ$ is expected to significantly diminish the fluxes of TDE flares due to their soft X-ray spectra. Second, the increasing density of foreground Galactic sources (primarily cataclysmic variables) can cause significant contamination of the sample. However, the contribution of Galactic contaminant sources remains noticeable (~50% of our raw candidate sample, see below) even for $|b| > 30^\circ$. In addition, we exclude sources with the angular extent larger than 30 arcsec, which allows us to avoid dealing with extended sources like galaxy clusters. Finally, we produce a list of positions that can be used as an input for the XMM flux upper limit server.

At the second stage, the XMM flux upper limit server reads the file produced previously and returns a FITS-file with fluxes in various energy bands measured within a circle centred on the XMM position of a given source, provided there are XMM-Newton observations of the corresponding field. We use the 0.2–2 keV energy band and an extraction radius of 60 arcsec, which exceeds the RASS localisation uncertainty, to obtain a conservative estimate. It is important to note that such a flux estimate will obviously include contributions from the background emission and from any sources falling into the extraction region. At energies above 1 keV, it is dominated by the nearly constant contribution of extragalactic sources with the total surface brightness of $4 \times 10^{-15}$ erg s $^{-1}$ cm$^{-2}$ arcmin$^{-2}$, while below 1 keV it is dominated by diffuse Galactic and local thermal-like emission, which varies from position to position and has the net surface brightness of the same order (Hickox & Markevitch 2006).

Hence, we can miss potential TDE candidates located in highly crowded regions or regions polluted by extended sources. In addition, we remove XMM-Newton observations with very small exposure time and hence with a high uncertainty of the flux estimate.

Finally, we check whether the flux derived from the XMM-Newton observations is at least 10 times lower than the lowest of flux estimates provided by RASS-BSC (i.e. Flux1 or Flux2). This resulted in a preliminary sample of 24 TDE candidates, selected based on X-ray variability properties only. Since there are only ~20 photons detected by ROSAT for each of these objects, we do not impose any additional constraints on our sample based on X-ray spectral properties of the sources (reflected in their hardness ratios, see Section 2.1).

We then additionally examined these candidate sources for any problems that might have significantly affected their positions and fluxes reported in RASS-BSC. To this end, we first checked for the presence of cautionary flags in RASS-BSC, which may reflect some difficulties met by the detection algorithm. In addition, we visually compared available RASS and XMM-Newton X-ray images for our candidates.

As a result, one candidate (1RXS J132654.5-271104), which is reported as displaying a complex diffuse emission pattern (‘d’-flag) but nevertheless considered point-like in RASS-BSC, is likely associated with the Abell 1736 galaxy cluster. The flux reported in RASS-BSC is collected from a much larger region (roughly, from the extraction region with a radius of 5 arcmin) compared to the region (with a radius of 1 arcmin) we used to get an upper limit from XMM-Newton and can be dominated by X-ray emission from the intracluster medium.

A somewhat different situation takes place in two other cases, 1RXS J003406.7-020935 and 1RXS J012605.2-012151, where the regions of interest appear to be contaminated by both a relatively large number of point sources and relatively weak diffuse emission from RXC J0034.6-208 (a galaxy cluster candidate from Böhringer et al. 2004) and Abell 194, respectively, as is clearly seen on XMM-Newton images (see also Hudaverdi et al. 2006 reporting the results of XMM-Newton data analysis for Abell 194). In both cases, one of the brightest point sources in the field falls rather close to the centroid provided by RASS-BSC, albeit well outside the cited localisation uncertainty region. Namely, it is 3XMM J003410.7-021039 (IC 0029), with a 0.2–12 keV flux of $2.8 \pm 0.5 \times 10^{-13}$ erg/s/cm$^2$, 88 arcseconds away from the RASS-BSC centroid for 1RXS J003406.7-020935 (having 1σ localisation error of 15 arcseconds), and 3XMM J012600.6-012041 (the NGC547/NGC545 pair), with a 0.2–12 keV flux of $6.0 \pm 0.2 \times 10^{-13}$ erg/s/cm$^2$, 98 arcseconds apart from the RASS-BSC centroid for 1RXS J012605.2-012151 (having 1σ localisation error of 25 arcseconds). Assuming that it is these sources which are in fact responsible for the fluxes reported in RASS-BSC, these candidates then fail to satisfy our TDE selection criterion based on high-amplitude variability.

Thus, there remain 21 candidates in our sample. However, some of them still have cautionary RASS-BSC flags due to the presence of a nearby bright source, and we will discuss this specifically in every interesting case below.

### 2.4 Cross-correlation with other surveys

Since many of the RASS-BSC sources have been identified and classified since they were discovered, we searched NASA/IPAC Extragalactic Database (NED) and the SIMBAD Astronomical Database for possible identifications of our TDE candidates. Nine of the candidates turned out to be known stellartype objects, including 7 systems with accreting white dwarfs (RX J0132.7-6554, Burwitz et al. 1997; EF, Eri, Verbunt et al. 1997, RX J0527.8-6954, Trümper et al. 1991; Greiner et al. 1996a, RX J1039.7-0507, Appenzeller et al. 1998, WGA J1047.1+6355, Singh et al. 1995, RX J1957.1-5738, Thomas et al. 1996, RX J2022.6-3954, Burwitz et al. 1997, a symbiotic star, CD-43°14304 Muerset, Wolff & Jordan 1997), and a hot white dwarf, WD 1232-546 (Pounds et al. 1993).

Four other candidates have been reported to demonstrate signatures of the active galactic nucleus (AGN) activity: KUG 1624+351 – a Seyfert 1.5 (Sy1.5) galaxy at $z = 0.0342$ (Bade et al. 1998, Véron-Cetty & Véron 2006), RX J1225.7+2055 – a narrow line Seyfert 1 (NLSy1) galaxy at $z = 0.335$ (Greiner et al. 1996b), MCG -01-05-031 – a Seyfert 2 (Sy2) galaxy at $z = 0.0182$ (Panessa & Bazzari 2002), and WPVS 007 – a NLSy1 galaxy at $z = 0.028$. The last object is famous for very unusual X-ray properties, primarily soft X-ray variability by two orders of magnitude over three years (Grupe et al. 1995) and has already been mentioned in the context of TDE searches (Donley et al. 2003). However, the very recent extensive monitoring of WPVS 007 by the Swift observatory clearly indicates a high level of short-term variability of the source consistent with a partial covering absorber model (Grupe et al. 2013). Having been observed by XMM-Newton for ~ 100 kiloseconds (2010 June 11), the source was identified as 3XMM J003915.8-511701 with the mean 0.2–12 keV flux of

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5 http://ned.ipac.caltech.edu/

6 http://simbad.u-strasbg.fr/simbad/
(1.12 ± 0.16) × 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2}, \text{ i.e.} \sim 500 \text{ times fainter compared with the detection during the RASS. The other three sources have also been marginally detected by XMM-Newton and Table 1 provides brief information about these detections.}

The remaining eight candidates (see Table 2 and Table 3) either do not have firm identification or are associated with galaxies with no obvious AGN signatures. We have looked for potential counterparts of these objects inside their ROSAT localisation regions using the Guide Star Catalog II (GSC-II)7 based on the Digitized Sky Survey (DSS) images (Lasker et al. 2008), the Two Micron All Sky Survey (2MASS, Skrutskie et al. 2006)8, the Wide-field Infrared Survey Explorer (WISE) all-sky survey9 (Wright et al. 2010) and the Sloan Digital Sky Survey (SDSS10). The next section presents the results of this search.11

3 RESULTS

Table 2 summarises the RASS-BSC data for the potential TDE candidates. Table 3 provides X-ray flux limits for these sources obtained more than 10 years later with XMM-Newton. Since the typical RASS localisation uncertainty is r_{\text{RASS}} \sim 15 \text{ arcsec}, it is difficult to uniquely associate these TDE candidates with sources from IR, optical or high-angular-resolution X-ray surveys, and we should regard all sources (if any) falling into the RASS localisation region as possible counterparts. We use the 2σ (i.e. \sim 95%) confidence region (corresponding to a radius r_{\text{2σ}} \sim 1.6 r_{\text{RASS}} for the two-dimensional Gaussian distribution) for such analysis. Hence, only a few per cent of true counterparts can be missed.

Typically, there are several possible counterparts in the RASS localisation region. For example, this is the case if SDSS photometric data are available. The easiest possibility of proceeding with identification in such cases is checking for sources with AGN or cataclysmic variable (CV) signatures and thus rejecting a TDE origin. This can be done in several ways. First, a TDE flare taking place in the RASS epoch is expected to disappear by the time of the XMM-Newton observations, while the X-ray emission from the TDE host (inactive) galaxy is very unlikely to exceed the XMM-Newton detection limit. In contrast, there can well be a significant X-ray detection in 3XMM-DR4 (of course, this flux being at least 10 times lower than the flux measured by ROSAT) in the case of an AGN or CV. Second, any spectral information could of course be helpful but, unfortunately, an optical spectrum is available only for a small fraction of the possible counterparts. Finally, WISE infrared (IR) colors can also indicate an AGN origin, and such information is indeed available for many of the possible counterparts.

In addition, we can put a simple, physically motivated constraint on the expected optical brightness of a TDE host galaxy. Indeed, given the mass of the central SMBH, M_{BH}, the (near) peak soft X-ray luminosity of a TDE is expected to be

\[ L_x = k L_{\text{Edd}}(M_{BH}) = kL_{\text{Edd,0}} M_{BH} / M_{\odot}, \quad (1) \]

where \( L_{\text{Edd,0}} \approx 1.4 \times 10^{38} \text{ erg/s} \) is the Eddington luminosity for one solar mass and the factor \( k \sim 0.1 \) accounts for the bolometric correction, possible geometrical dilution and the fact that the X-ray flux detected by RASS may be measured not exactly at the peak of the TDE (see e.g. Khabibullin, Sazonov, & Sunyaev 2014).

On the other hand, the optical luminosity of the TDE host galaxy could be roughly expressed as

\[ L_V = \frac{L_{\text{Edd}}}{P_b} = \frac{200}{P_b} \left( \frac{M_{BH}}{M_{\odot}} \right)^{1.2} \quad (2) \]

where \( L_{\odot} \approx 4.6 \times 10^{33} \text{ erg/s} \) is the solar V-band luminosity, \( P_b \sim 0.1 \) is the fraction of the bulge in the total galaxy luminosity, and we have converted \( L_{\text{Edd}} \) to \( M_{BH} \) using the \( M_{BH} - L \) relation from Gilletkin et al. (2009) for \( M_{BH} \sim 10^6 M_\odot \) (in fact, the \( M_{BH} - L \) relation is almost linear, so the numerical factor in Eq. 2 depends on \( M_{BH} \) very weakly).

As a result, we obtain the following relation between the TDE soft X-ray \((-0.2–2 \text{ keV})\) flux, \( F_x \), and the optical flux of the host galaxy, \( F_V \):

\[ F_V = \frac{F_x L_x}{k L_{\text{Edd,0}}} \approx \frac{200}{k P_b} \left( \frac{M_{BH}}{M_{\odot}} \right)^{0.64} F_x / P_b, \quad (3) \]

where \( k = k/0.1 \) and \( P_b = P_b/0.1 \).

For \( F_x = 3 \times 10^{-13} \text{ erg/s/cm}^2 \), i.e. close to the RASS-BSC detection limit, \( F_V = 1.9 \times 10^{-14} \text{ erg/s/cm}^2 \), given \( k = 0.1 \) and \( P_b = 0.1 \), which translates to \( m_V = 2.5(-5.5 - \log F_x) \approx 20.55 \). For \( F_x = 1 \times 10^{-12} \text{ erg/s/cm}^2 \), a typical flux for RASS-BSC sources (and for our candidate sources, see Table 2), one gets \( m_V \approx 19.25 \).

Thus, a TDE host galaxy is unlikely to be very faint, unless we deal with an outlier from the \( M_{BH} - L \) relation. We caution though that the \( M_{BH} - L \) correlation used above was derived from a sample with somewhat heavier SMBHs compared to those expected to produce most TDEs in the Universe.

Similar estimates can be done for the near-IR (e.g. K band) magnitude of TDE host galaxies. According to the \( M_{BH} - L_k \) relation for the total (bulge+disk) K-band luminosity from Lisker et al. (2014),

\[ L_K \approx 300 L_{\odot} \frac{M_{BH}}{M_\odot} \quad (4) \]

for \( M_{BH} \sim 10^6 M_\odot \), where \( L_{\odot} \approx 8 \times 10^{30} \text{ erg/s} \) is the solar K-band luminosity (e.g. Binney & Merrifield 1998), which corresponds to a Vega magnitude \( M_K \approx 3.28 \). Therefore,

\[ F_K = \frac{F_x 300 L_{\odot}}{k L_{\text{Edd,0}}} \approx 2 \times 10^{-5} F_x / k \quad (5) \]

Hence, for \( F_x = 1 \times 10^{-12} \text{ erg/s/cm}^2 \), \( F_K = 2 \times 10^{-15} \text{ erg/s/cm}^2 \), or \( m_K \approx 19.4 \).

For different classes of contaminating sources, yet simpler calculations can be carried out. As known from X-ray surveys, 0.1 \( \leq F_x / F_V \leq 10 \) for AGN and CVs in the soft X-ray band (e.g. Maccacaro et al. 1988; Aird et al. 2010), which translates to 15-15 \( \leq m_K \leq 20.15 \) given \( F_x = 3 \times 10^{-13} \text{ erg/s/cm}^2 \). Some higher \( F_x / F_V \) ratios could be found during flares of luminous blazars (Maccacaro et al. 1988; Cesco et al. 2012). It should

\(^{12}\) This value for \( P_b \) is typical for disk galaxies but not for elliptical ones. However, for elliptical galaxies \( M_{BH} \) typically exceeds the tidal disruption threshold for solar-type stars \((-10^8 M_\odot)\), so we do not expect significant contribution from such galaxies in our TDE sample.

\(^{13}\) In this case, the \( M_{BH} - L_k \) relation is also almost linear, so the numerical factor in Eq. 4 will be essentially the same for \( M_{BH} \sim 10^6 M_\odot \).

\(^{14}\) The flux-to-magnitude conversion was calculated as \( m_K = 2.5(-6.95 - \log F_K) \), in accordance with the 2MASS photometric system (Cohen, Wheaton, & Megeath 2003).
be mentioned however that in the case of AGN and CVs, optical
counterpart emission can be variable and tightly related to the X-
ray emission. So if we compare the (relatively high) X-ray flux
measured during the RASS with non-contemporaneous optical ob-
servations (e.g. SDSS), the resulting X-ray/optical flux ratio may
be different, likely higher, than estimated above, hence the optical
counterpart may be somewhat fainter.

For flaring stars, $F_X/F_V \sim 10^{−3}−10^{−2}$ is much more typical
(with rare exceptions of very intense flares in low-mass stars,
Favata & Micela 2003), so the optical counterpart is expected to be
brighter than $m_V \approx 13$ in this case.

In summary, the optical counterparts of our candidate sources are
likely to be present in the SDSS photometric data (complete to $r \sim 22.2$
Abazajian et al. 2004), if available, even in the case of AGN/CV association, and they are probably significantly bright to be present in DSS images in the case of TDE association (GSC-II is complete to $R_p \approx 20$, Lasker et al. 2008).

### 3.1 1RXS J002048.5-253823

1RXS J002048.5-253823 is located approximately 5 arcminutes
from the galaxy cluster Abell 022, which is present in RASS-
BSC as an independent X-ray source, 1RXS J002041.8-253437, and is
~6 times brighter than 1RXS J002048.5-253823. Therefore, the
localisation and the flux estimate could be somewhat biased
for 1RXS J002048.5-253823. However, examination of the RASS
image shows that the two sources are well separated from each other. The available crude spectral information (hardness ratios), $HR1\lesssim −0.32 \pm 0.27$ and $HR2\lesssim −0.82 \pm 0.87$ for 1RXS J002048.5-253823 against HR1$\approx −0.31 \pm 0.14$ and HR2$\approx −0.22 \pm 0.15$ for 1RXS J002041.8-253437, is of little value in this respect given the large uncertainties for the fainter source.

Moreover, there is a faint 3XMM-DR4 source (3XMM J002049.3-253828) with the 0.2–2 keV flux of $(1.7 \pm 0.3) \times 10^{−14}$
ergs/cm², i.e. close to the detection limit of the 3XMM-DR4 cata-
logue, inside the 1σ-region (12 arcsec away from the reported
RASS-BSC centroid). This source has an optical counterpart in the DSS image (and the corresponding entry in GSC-II with $R_p = 17.1$), a near-infrared counterpart in the 2MASS catalogue and a mid-infrared counterpart in the WISE catalogue (see Fig. 1 and Table 4). The WISE colors (namely, $W2 − W1 = 1.24 \pm 0.18$) strongly indicate an AGN association for this source (Stern et al. 2012).

Since there are no other objects inside the RASS localisation re-
region, we suggest that 1RXS J002048.5-253823 is a high-amplitude

### Table 1. Strongly variable X-ray sources with AGN signatures

| Name       | Type$^a$ | $z^a$ | 3XMM ID     | 3XMM-DR4 Date (UTC) | Flux$^b$ | RASS-BSC Count rate$^c$ | Flux1$^c$ | Flux2$^c$ | Amplitude$^d$ |
|------------|---------|-------|-------------|---------------------|---------|------------------------|----------|----------|-------------|
| WPVS 007   | NLSty1  | 0.028 | J003915.8−511701 | 2010-06-11.9        | 1.12 ± 0.16 | 0.96 ± 0.068           | 19.4     | 3.04     | > 270       |
| MCG-01-05-031 | Sy2   | 0.0182 | J014525.4−034938 | 2008-08-03.2        | 5.97 ± 0.7 | 0.17 ± 0.03            | 3.5      | 0.92     | > 15        |
| RX J1225.7+2055 | NLSy1 | 0.335 | J122541.9+205503 | 2003-06-12.7        | 9.95 ± 1.19 | 0.33 ± 0.03            | 6.7      | 1.63     | > 16        |
| KUG 1624+351 | Sy1.5  | 0.0342 | J162636.5+350241 | 2007-08-17.1        | 3.68 ± 0.708 | 0.081 ± 0.013        | 0.971    | 0.672    | > 18        |

$^a$ See references in the text.
$^b$ The 0.2–12 keV flux measured by XMM-Newton, $10^{−14}$ erg s$^{−1}$ cm$^{−2}$.
$^c$ ROSAT count rate with a 1σ uncertainty (cts/s) and corresponding estimates for the 0.1–2.4 keV flux (Flux1 and Flux2, see text), $10^{−12}$ erg s$^{−1}$ cm$^{−2}$.

### Table 2. RASS-BSC data for potential TDE candidates

| 1RXS ID    | RA          | Dec         | PosErr$^d$ | CR$^b$ | eCR$^b$ | Exp$^b$ | HR1$^c$ | eHR1$^c$ | HR2$^c$ | eHR2$^c$ | Flux1$^d$ | Flux2$^d$ | $N_e^{1}$ | $nH$ |
|------------|-------------|-------------|------------|--------|---------|---------|---------|----------|---------|----------|----------|----------|----------|------|
| J002048.5−253823 | 5.20208 | −25.63986 | 16(26) | 0.0563 | 0.0167 | 280     | −0.32  | 0.27     | −0.82   | 0.87     | 9.78     | 3.72     | 2.3 |
| J005626.3−010615 | 14.10958 | −1.10431 | 15(24) | 0.0999 | 0.0202 | 312     | 0.65   | 0.18     | 0.61    | 0.19     | 21.5     | 11.7     | 3.4 |
| J112312.7+012558 | 170.80292 | 1.48292 | 17(28) | 0.052 | 0.0143 | 375     | 0.08   | 0.27     | 0.15    | 0.35     | 13.7     | 4.53     | 3.5 |
| J114727.1+494302 | 176.86292 | 49.71722 | 9(15) | 0.414 | 0.0402 | 299     | −0.97  | 0.02     | 1.13    | 1.3      | 63.3     | 13.1     | 1.7 |
| J105447.2+641225 | 196.44666 | 64.21445 | 9(15) | 0.167 | 0.0207 | 544     | −0.82  | 0.06     | −0.62   | 0.28     | 23.2     | 6.61     | 1.5 |
| J215101.5−302852 | 327.75626 | −30.48111 | 14(23) | 0.0534 | 0.0159 | 333     | −0.34  | 0.24     | 0.42    | 0.47     | 8.29     | 3.47     | 1.7 |
| J235424.5−102053 | 358.60208 | −10.34806 | 24(39) | 0.0523 | 0.0153 | 310     | 1      | 0.25     | −0.13   | 0.29     | 10.8     | 7.12     | 2.7 |

$^d$ 1σ localisation uncertainty and corresponding radius of the 95% confidence region, arcsec.
$^e$ The source mean count rate and corresponding 1σ error (both in units of counts s$^{−1}$) and exposure time (in sec).
$^f$ Hardness ratios (see text) with corresponding 1σ errors.
$^g$ 0.2–2.4 keV flux estimates (see text), $10^{−13}$ erg s$^{−1}$ cm$^{−2}$.
$^h$ Total Galactic H I Column Density (divided by $10^{20}$ cm$^{−2}$) as determined with the nH utility of the FTOOLS package (based on the map of Kalberla et al. 2005).
The Digitized Sky Survey (DSS) image (2’x2’) of the region around 1RXS J002048.5-253823. The dashed and solid blue circles contain the 1σ and 2σ regions around the centroid of the ROSAT source, respectively. Red squares mark 2MASS sources, green crosses mark 3XMM sources (if any). For those fields with SDSS data available (see the subsequent figures), black diamonds mark the positions of SDSS sources. The arrow points out the most probable counterpart.

(≈ 50 if Flux1 is used as a flux estimate for the RASS-epoch) AGN flare.

We conclude that 1RXS J002048.5-253823 is probably not a TDE.

3.2 1RXS J005626.3-010615

This candidate is also located nearby a galaxy cluster, Abell 119, though approximately 10 arcminutes from its centre. Abell 199 enters RASS-BSC as 1RXS J005617.5-011501, being ~5 times brighter than 1RXS J005626.3-010615 with the extension of ≥2 arcmin. So, the sources seem to be well separated for the extraction region of the fainter one (10 arcmin in diameter) to include almost no extended emission from the brighter one. Both hardness ratios are also formally different for them: HR1=0.65 ± 0.18, HR2=0.61 ± 0.19 for 1RXS J005626.3-010615 against HR1=0.88 ± 0.05, HR2=0.29 ± 0.08 for 1RXS J005617.5-011501.

There are seven SDSS sources inside the localisation region (see Fig. 3). All of them are faint, r ≥ 20, and thus invisible in the DSS image. An optical spectrum is available for only one source, which can be unambiguously classified as an M-star. This source (24 arcsec from the RASS centroid) is also the only one with a 2MASS counterpart. Although such a star could potentially produce an intense X-ray flare, it should be much brighter since the X-ray flux in the ROSAT energy range exceeded 10^{12} erg/s/cm^2 and one should then expect m_v ≤ 16.5 assuming F_x/F_v ≤ 1 (such ratios correspond to the most extreme stellar flares).

Three other sources are identified as extended ones and another three are consistent with being point-like. The data on these sources are given in Table 4. An extended source 16 arcsec away from the RASS-BSC centroid is the only one detected by WISE and its colors favour an AGN association. It is also the optically brightest one with r ≈ 21.0. This object can thus be considered the most probable counterpart. As to its nature, it could be a highly variable AGN (perhaps a blazar) with a very high F_x/F_v ratio.

We thus conclude that 1RXS J005626.3-010615 is unlikely to be a TDE.

3.3 1RXS J101326.2+061202

There are two SDSS sources at the boundary of the localisation region, 25 arcsec away from the RASS-BSC centroid (see Fig. 4). One is classified as a K star according to the SDSS spectrum, so it potentially could be responsible for the highly variable X-ray emission. Indeed, the optical magnitudes of the source satisfy the condition discussed above. This source is also the only one inside the uncertainty region with 2MASS (with JHK magnitudes of 15.56, 14.86 and 14.82) and WISE (with W1 = 14.69 and...
Tidal disruption events in RASS-BSC

Figure 2. DSS image (2'x2') of the region around 1RXS J005626.3-010615. The symbols are the same as for Fig. 1.

Figure 2

Figure 3. DSS image (2'x2') of the region around 1RXS J101326.2+061202. The symbols are the same as for Fig. 1.

Figure 3

W2 = 14.79) counterparts. The other source is too faint to even be considered a possible counterpart in the case of a TDE association (see Table 4). Thus, we suggest that the most probable counterpart is the K star, which experienced a strong flare during the RASS.

We conclude that 1RXS J101326.2+061202 is unlikely to be a TDE.

3.4 1RXS J112312.7+012858

There is a 3XMM-DR4 source (3XMM J112313.1+012924), with a flux of \((6.0 \pm 0.8) \times 10^{-14} \text{ erg/s/cm}^2\) in the 0.2–2 keV energy band, at the boundary of the localisation region (28 arcsec from the RASS-BSC centroid). This source has a faint, non-extended SDSS counterpart with \((22.2, 21.3, 21.0, 21.4, 20.8) ugriz\) magnitudes, but there are no 2MASS or WISE counterparts. Taking these facts together, we consider association of the 3XMM-DR4 source with 1RXS J112312.7+012858 unlikely.

In total, there are eleven SDSS sources inside the localisation region. One of these has been confidently identified as an A0 star by means of optical spectroscopy. The data on all other sources are summarized in Table 4. The brightest SDSS object \((r = 19.9)\), 21 arcsec away from the RASS-BSC centroid, is consistent with an extended source and is also the only one with a WISE counterpart. The first WISE color \((W1 - W2 = 15.114 - 14.831 = 0.283)\) makes a luminous AGN association unlikely for this source. However, this object can be the host galaxy of a TDE.

We conclude that 1RXS J112312.7+012858 is possibly a TDE.

3.5 1RXS J114727.1+494302

This source, also known as RBS 1032, has been intensively investigated by Ghosh et al. (2006). They reported an ultrasoft X-ray spectrum, consistent with a blackbody with \(kT_{bb} \approx 70\) eV, as well as a \(\sim 3\)-fold flux drop from the original RASS detection (November 05, 1990) to the first pointed ROSAT observation carried out \(\sim 2.1\) years later (December 07, 1992) and a \(\sim 6\)-fold drop to the second pointed observation \(\sim 3.6\) years later (June 05, 1994) (Ghosh et al. 2006). The authors identified RBS 1032 with a dwarf spheroidal galaxy (SDSS J114726.69+494257.8) at \(z = 0.026\) with no optical signatures of nuclear activity. The WISE colours of this source
(W1 – W2 = 0.115) also rule out a luminous AGN. XMM-Newton detected only a marginal signal from this position in late 2009 with the net 0.2–12 keV flux of (1.1 ± 0.6) × 10^{-14} erg s^{-1} cm^{-2}.

There are two other SDSS extended sources inside the localisation region, with neither 2MASS nor WISE detection (see Fig. 6 and Table 4). These sources are too faint to be considered potential TDE hosts.

Ghosh et al. (2006) argued that the variable X-ray emission may come from an intermediate mass black hole accreting matter from a white dwarf companion. The estimated soft X-ray luminosity at peak is ~ 10^{41} ergs/cm^2 and fitting the X-ray spectra with a blackbody model yields a temperature of kT_{bb} ≈ 70 eV (Ghosh et al. 2006). Both of these characteristics are consistent with the TDE scenario (Komossa 2002; Esquej et al. 2008). However, the six-fold flux decrease in the 3.6 years spanned by the ROSAT observations may seem to be too small for the TDE scenario assuming that the first detection occurred near the TDE peak.

However, the estimated luminosity is not sufficiently high to exclude the possibility that the first detection in fact occurred at a relatively late phase of the event. Indeed, assuming the canonical shape of the decay phase of the TDE light curve (Phinney 1989),

\[
\frac{L(\tau + \delta t)}{L(\tau)} = \left(\frac{\tau + \delta t}{\tau}\right)^{5/3},
\]

where \(\tau\) is the time of the RASS observation relative to the disruption moment and \(\delta t = 3.6\) yrs is the time passed between the RASS observation and the second pointed ROSAT observation. Given that the unabsorbed 0.1–2 keV flux decreased from 6.0 × 10^{-12} erg s^{-1} cm^{-2} to 1.1 × 10^{-12} erg s^{-1} cm^{-2} over this period (Ghosh et al. 2006), we find \(L(\tau + \delta t)/L(\tau) = 1.1/60\), so \(\tau = 2.0\) yrs. This implies that the disruption took place near the 5th of November, 1988. Since the duration of the early (Eddington-limited) phase \(\tau_{Edd} \approx 0.25 (M_{BH}/10^7 M_\odot)^{-5/3}\) for the minimal peri-center distance, \(R_p = 3R_s\) (with \(R_s = 2GM_{BH}/c^2\) being the Schwarzschild radius of the black hole), of the disrupted star (e.g. Strubbe & Quataert 2009), the peak luminosity is expected to be as high as \(L_{peak} \approx 30L(\tau) \approx 3 \times 10^{46}\) erg/s for \(M_{BH} = 10^7 M_\odot\), which is nonetheless ~5 times smaller than the Eddington luminosity for this black hole mass. The same calculation for \(M_{BH} = 5 \times 10^6 M_\odot\) (\(\tau_{Edd} \approx 0.19\) yr) results in \(L_{peak} \approx 0.5L(\tau) \approx 5 \times 10^{44}\) erg/s, i.e. just slightly below the corresponding Eddington luminosity. Although these estimates may be slightly affected by the bolometric correction, which can be estimated as \(L_{bol} \approx 1.1L_{X,0.1-2.4}\) from the observed shape of the spectrum (a black body with kT_{bb} = 70 eV), and geometrical effects, it appears that values \(M_{BH} \lesssim 5 \times 10^6 M_\odot\) are unlikely.

Therefore, the X-ray light curve seems to be consistent with a TDE that occurred near a ‘normal’ SMBH with \(M_{BH} = 5 \times 10^6 – 10^7 M_\odot\), even though a simple power-law model predicts an order of magnitude higher flux for the XMM-Newton epoch (see Fig. 6 for an illustration) – this is in fact not surprising since at this very late stage of the TDE flare the feeding rate of the SMBH is expected to have fallen by several orders of magnitude.

We note that for \(M_{BH} = 5 \times 10^6 – 10^7 M_\odot\), one could expect a somewhat lower temperature of the multicolour blackbody accretion disk, kT_{bb} \approx 40 – 50 eV \propto M_{BH}^{-1/4} (Shakura & Sunyaev 1973). However, at late phases (as suggested here for the ROSAT observations) when the luminosity is as low as a few percent of the peak value, i.e. the mass accretion rate is a few per cent of the critical accretion rate, the X-ray spectrum may harden due to transition to an inefficiently cooling accretion flow and/or formation of relativistic jets.

Therefore, 1RXS J114727.1+494302=RBS 1032 is probably a tidal disruption event associated with a SMBH of mass \(M_{BH} \sim 5 \times 10^6 – 10^7 M_\odot\) which was caught by ROSAT at a fairly late phase. We note however that the inferred black-hole mass may be too high for the dwarf host galaxy of RBS 1032 (Ghosh et al. 2006).
3.6 1RXS J130547.2+641252

There is only one SDSS source inside the localisation region (3 arcsec away from the RASS-BSC centroid, see Fig. 7). It is consistent with an extended source and has (22.5, 22.0, 20.9, 20.6, 20.4) ugriz magnitudes. This object is fainter than would be expected for a TDE host galaxy, albeit not very dramatically.

We conclude that 1RXS J130547.2+641252 is possibly a TDE.

3.7 1RXS J215101.5-302852

A weak 3XMM-DR4 source (3XMM J215100.7-302832), \((4.1 \pm 0.8) \times 10^{-14} \text{ erg/cm}^2 \) (0.2–2 keV), is detected at the boundary of the localisation region (22 arcsec from the RASS-BSC centroid). This source has an optical counterpart in the DSS image (with \( R = 18.6 \)) as well as infrared counterparts in the 2MASS (with \( JHK \) magnitudes of 16.849, 16.304 and 15.347) and WISE (with \( W1, W2, W3 \) magnitudes of 13.985, 12.949 and 10.192) surveys (see Fig. 8). The first WISE colour, \( W1 - W2 = 1.036 \), strongly indicates an AGN origin.

We conclude that 1RXS J215101.5-302852 is unlikely to be a TDE.

3.8 1RXS J235424.5-102053

1RXS J235424.5-102053 is located approximately 5 arcminutes from the centre of the galaxy cluster Abell 2670. The latter enters RASS-BSC as 1RXS J235409.4-102506, is \( \sim 10 \) times brighter than 1RXS J235424.5-102053 and has an angular extent of \( \sim 1.5 \) arcmin. The uncertainties in the hardness ratios for 1RXS J235424.5-102053 are too large to draw any conclusions based on them. We note however that the corresponding spectrum should be rather peculiar with HR1=1.0±0.25, i.e. having no counts in the 0.1–0.4 keV energy channel, and HR2=0.13±0.29 (while 1RXS J235409.4-102506 has HR1=0.83 ± 0.06 and HR2=0.24 ± 0.1).

The most prominent candidate for the counterpart of this source is a bright SDSS galaxy (17 arcsec away from the RASS-BSC centroid, see Fig. 9) at \( z = 0.0805 \), which may be a distant member of the Abell 2670 cluster (\( z = 0.076 \)). Its optical spectrum does not show signatures of luminous nuclear activity, which is confirmed by the colours of the corresponding WISE counterpart (\( W1 - W2 = 0.0 \)). All other eight SDSS sources inside the localisation region are faint (\( r > 20 \)) and not detected by WISE. We further note that the bright SDSS galaxy appears to be a disk one observed nearly edge-on, which might cause significant absorption of the X-ray emission in the softest energy band of ROSAT and thus explain the measured value of HR1 close to unity.

We conclude that 1RXS J235424.5-102053 could still be considered a possible TDE candidate. However, we cannot exclude that it is a spurious point-like source on a RASS image contaminated by extended cluster emission.

4 DISCUSSION

Since any TDEs that may be present in our sample occurred almost 25 years ago, there is not much that can be done to explore their individual properties (see e.g. Bower et al. 2013 who looked for the late-time radio emission from TDE candidates discovered during the RASS). Nonetheless, this sample is valuable for studying TDE population properties. In particular, it allows us to put a constraint on the rate of TDEs in the local Universe.

4.1 Comparison with the results of Donley et al.

We first compare our results with the findings of Donley et al. (2002). Their sample was limited to relatively bright X-ray sources, with the unabsorbed 0.2–2.4 keV flux during the RASS higher than \( f_0 = 2 \times 10^{-12} \text{ erg/cm}^2 \) (assuming a power-law spectrum with \( \Gamma = 4 \)), which was due to the sensitivity threshold of the subsequent ROSAT pointed observations (Donley et al. 2002) used in the
Table 4. Possible counterparts of potential TDE candidates. The coordinates of the most probable counterparts are marked in boldface.

| Candidate | RA   | Dec  | Dist | WISE, W1–W2 | 2MASS, JHKs | SDSS, Class (ugriz) | 3XMM | Nature   |
|-----------|------|------|------|-------------|-------------|---------------------|-------|----------|
| 1RXS J002048.5-253823 | 5.2054 | -25.6414 | 12 | 13.3–12.1 | 16.5–16.1–15.0 | s(17,33,17,32,17,74) | 1.68 ± 0.46 | AGN |
| 14.1057 | -1.1045 | 13 | 1RXS J005626.3-010615 | 14.1066 | -1.1075 | 16 | 16.8–15.7 | s(25,03,25,6,22,3,22,5) | (22,52,22,17,22,0,21,0) | AGN |
| 14.1054 | -1.1067 | 20 | 14.1093 | -1.1106 | 23 | 14.1146 | -1.1100 | 24 | 14.8–14.6 | 15.8–15.1–14.9 | s(22,6,1,8,20,1,18,3,17,3) | (23,9,2,4,2,2,5,2,1,4,20,9) | M-star |
| 14.1110 | -1.1108 | 24 | 153.3579 | 6.1939 | 25 | 14.7–14.8 | 15.6–14.9–14.8 | s(21,4,1,8,17,4,17,0,16,7) | (24,6,2,4,2,1,24,2,2,1,2,2,1) | K-star |
| 135.3523 | 6.2021 | 25 | 170.8046 | 1.4824 | 7 | 170.8063 | 1.4845 | 13 | 170.8063 | 1.4859 | 16 | 170.8065 | 1.4783 | 21 | 15.1–14.8 | s(24,1,2,1,6,19,9,19,2,1,8) | (22,4,2,2,1,7,21,2,4,2,1,1) | TDE |
| 170.8082 | 1.4788 | 24 | 170.8084 | 1.4781 | 26 | 170.7954 | 1.4818 | 27 | 170.8050 | 1.4903 | 28 | 170.8050 | 1.4901 | 28 | 170.8065 | 1.4903 | 28 | 24.1,2,1,6,21,1,4,2,0,2 | 6.0 ± 0.8 | AGN/CV? |
| 176.8613 | 49.7161 | 6 | 176.8618 | 49.7143 | 11 | 176.8656 | 49.7133 | 15 | 196.4486 | 64.2145 | 3 | 176.8613 | 49.7161 | 6 | 14.5–14.3 | 15.9–15.1–15.2 | g(19,1,1,7,1,1,1,1,6,1,6,1,6) | (23,9,2,7,7,2,7,2,3,2,2,2,0) | TDE (z = 0.026) |
| 1RXS J114727.1+494302 | 1RXS J130547.2+641252 | 1RXS J215101.5-302852 | 327.7528 | -30.4758 | 22 | 14.0–13.0 | 16.8–16.3–15.3 | g(18,3,4,15,5,17,7) | 4.1 ± 0.8 | AGN |
| 358.6027 | -10.3484 | 2 | 358.6047 | -10.3501 | 12 | 358.6017 | -10.3446 | 13 | 358.5976 | -10.3467 | 17 | 13.7–13.7 | 15.4–14.7–14.3 | g(19,6,1,7,1,6,1,6,1,4,1,6) | (23,1,2,1,8,20,8,20,8,20,8) | TDE (z = 0.081) |
| 358.5976 | -10.3467 | 17 | 358.6036 | -10.3538 | 21 | 358.5976 | -10.3467 | 28 | 358.5976 | -10.3467 | 29 | 358.5976 | -10.3467 | 33 | 17.0–16.5 | g(23,7,2,2,2,2,2,2,2,2,2,2,2,2) | (23,7,2,2,2,2,2,2,2,2,2,2,2,2) | AGN |
| 358.6003 | -10.3578 | 36 | 16.8–16.8 | g(23,5,2,2,4,2,2,2,1,2,2,0,8) | (23,5,2,2,4,2,2,2,1,2,2,0,8) | AGN |

Notes:

- Distance from the centroid of the RASS-BSC localisation, arcsec.
- WISE photometry becomes unreliable at approximately W1 = 14, W2 = 13.5. Values W3 > 12 (as is typical of the detections under consideration) should be regarded as marginal detections, hence we make no use of the W3 magnitude in our analysis (see text).
- A source type (s for stellar-like and g for galaxy-like) is given along with the agriz magnitudes.
- The X-ray (0.2–2 keV) flux of the source in the 3XMM-DR4 catalogue, 10⁻¹⁴ erg s⁻¹ cm⁻².
- DSS data is provided if no SDSS data is available: type (s for point-like and g for an extended source) and (B, R, I) magnitudes, taken from the second generation Guide Star Catalog (Lasker et al. 2008). The characteristic uncertainty for these photometric magnitudes is ≈ 0.2 for sources at high (|b| > 30°) galactic latitudes (Lasker et al. 2008).

Analysis for demonstrating that the flux of a given RASS source has undergone a strong decline. Since we use the more sensitive XMM-Newton observations for the same purpose, our resulting sample goes all the way down to the sensitivity limit of the RASS-BSC catalogue of f = 3 × 10⁻¹³ ergs/cm² (see Table 2). As a result, we can estimate the expected number of TDE candidates in our sample from the number of candidates, N₀, found by Donley et al. (2002) as follows:

\[ N = N_0 \times \frac{S}{S_0} \times \left( \frac{f}{f_0} \right)^{-3/2} \tag{7} \]

where \( S/S_0 \) is the ratio of the sky areas covered by the two surveys and \( f/f_0 \) is the ratio of the corresponding limiting fluxes. This equation is approximate since it disregards reddening of the distant objects and assumes a simple power-law (with a slope of -3/2) log N—log S distribution, with both of these assumptions being somewhat inaccurate (see e.g. Khabibullin, Sazonov, & Sunyaev 2014).

The sky coverage of the XMM-Newton observations corresponding to the 3XMM-DR4 catalogue is 794 deg², i.e. almost 2% of the whole sky, which is 19th of the area covered by the ROSAT pointed observations (Donley et al. 2002). Given N₀ = 6 candidates (including SBS 1620+545 with a hard X-ray spectrum and two candidates with AGN signatures) found by Donley et al. (2002) at |b| > 30°, we find from the above equation N = 11.5, roughly in agreement with the number (eight) of TDE candidates
we actually found in our study (as well as with the number anticipated by Donley et al. 2002 for such a XMM-Newton-RASS cross-correlation study).

However, as we showed in Section 3 three of our candidates (1RXS J002048.5-253823, 1RXS J005626.3-010615 and 1RXS J215101.5-302852) are likely to be high-amplitude AGN flares, and yet another one (1RXS J101326.2+061202) might be associated with an intense flare from a K-star. Finally, three candidates are broadly consistent with the expectations of the TDE scenario. Namely, 1RXS J114727.1+494302, 1RXS J130547.2+641252 and 1RXS J235424.5-102053 can be associated with sufficiently bright galaxies without signatures of luminous AGN, although we cannot exclude that 1RXS J235424.5-102053 is a spurious RASS source. For the fourth source, 1RXS J112312.7+012858, a TDE association is also acceptable but an AGN association could not be rejected conclusively.

Thus, among our eight candidates, 3 or 4 could be associated with AGN, 1 with a flaring star and 2 to 4 with TDEs. Interestingly, Donley et al. (2002) found a similar proportion of candidates for AGN flares and TDEs, which supports the conclusion above that the numbers of TDEs found in these two studies are broadly consistent with each other.

4.2 TDE rate in the local Universe

The simple estimates presented above disregard a number of important points. First, by increasing the sensitivity of a survey, one gets the possibility not only to detect more distant TDE flares, but also flares at a later phase of their decay, i.e. ‘older’ ones at closer distances. Second, the problem in hand involves rather soft X-ray sources with the spectral maximum near or even below the ROSAT sensitivity range. This, along with the importance of the interstellar absorption in this energy range, makes the estimates quite sensitive to the cosmological reddening of the spectra. In addition, the non-cubic dependence of the comoving volume on the luminosity distance needs to be accounted for when dealing with redshifts $\sim 0.15$. Finally, the sky coverage of XMM-Newton observations is not completely independent of the positions of RASS-BSC sources, since some of them were the targets of XMM-Newton observations.

To estimate the last effect, we first find that $\sim 8\%$ of all RASS-BSC extragalactic ($|b| > 30^\circ$) sources have ever been observed by XMM-Newton, either on purpose or serendipitously, whereas the sky coverage of the 3XMM-DR4 catalogue is only $2\%$ of that for the RASS (essentially the whole sky). However, the relative fraction of XMM-Newton targets among RASS-BSC sources should depend on the class of object of interest. As concerns the present study, which is aimed at non-extended, variable sources, the only potentially important source of bias might be associated with clusters of galaxies, which have often been targets of observations for XMM-Newton. Indeed, since clusters contain a lot of galaxies, they are expected to be cites of strongly increased TDE activity. However, since our study probes the Universe out to $z \sim 0.2$, an XMM-Newton ($\sim 15^\prime \times 15^\prime$) field of view randomly located on the sky will contain more galaxies than a rich cluster, hence there should be no significant bias due to XMM-Newton pointing at galaxy clusters. Moreover, the fact that a significant part of the 3XMM-DR4 catalogue is based on observations of fields containing galaxy clusters may in fact result in a decrease in the effective coverage for our study, since ROSAT’s rather poor angular resolution prevents finding TDE flares in the vicinity of cluster cores.

Indeed, only two out of eight candidates in our final sample of TDE candidates, namely, 1RXS J114727.1+494302 and 1RXS J130547.2+641252, prove to have been targets of XMM-Newton observations, while the others were observed serendipitously. Thus, the XMM-Newton sky coverage can indeed be considered almost independent of RASS-BSC sources when dealing with TDE candidates.

Following the treatment in Khabibullin, Sazonov, & Sunyaev (2014), we assume now some characteristic shape for the light curve, spectrum and energetics of a TDE. Specifically, given the RASS-BSC sources may in fact result in a decrease in the effective coverage for our study, since ROSAT’s rather poor angular resolution prevents finding TDE flares in the vicinity of cluster cores.

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Following the treatment in Khabibullin, Sazonov, & Sunyaev (2014), we assume now some characteristic shape for the light curve, spectrum and energetics of a TDE. Specifically, given that the Galactic absorption in the rest frame of the TDE). The original spectrum will be modified by interstellar absorption both in the host galaxy and along the line-of-sight in our Galaxy. The median value of the Galactic HI column density for the sources in our sample is $N_H \approx 2.7 \times 10^{20}$ cm$^{-2}$ (see Table E). The absorption due to the gas in the host galaxy could be of the same order but is poorly known (see e.g. the comparison of the net $N_H$ value found by the fitting the X-ray spectra for one of our candidates, RBS J1032, to the corresponding Galactic absorption in that direction, Ghosh et al. 2004). Just for convenience, we adopt the host’s absorption column density $N_H^{host} \approx 2.3 \times 10^{20}$ cm$^{-2}$, so that the net column density $N_H \approx 5 \times 10^{20}$ cm$^{-2}$ (note that the Galactic absorption is imposed in the observer’s reference frame, while the host’s absorption in the rest frame of the TDE).

In order to calculate the limiting redshift for a TDE flare at which it could still be identified as a bright ROSAT source (i.e. with the count rate more than 0.05 cts/s), we performed a fake XSPEC simulation using XSPEC (Dorman & Arnaud 2001) and the ROSAT response matrix as provided by the NASA’s High Energy Astrophysics Science Research Archive Center (HEASARC). Assuming $M_{ bol} = 5 \times 10^6 M_\odot$, the limiting redshift turns out to be $z_{lim} \approx 0.18$.

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16 https://heasarc.gsfc.nasa.gov/docs/rosat/rosgof.html
This can be considered the characteristic depth of our sample. Using this value and repeating the calculation presented in Section 4.1 of (Khabibullin, Sazonov, & Sunyaev 2014), we finally can estimate the TDE rate in the local Universe at $R \approx 8.4 \times 10^{-3} \text{ Mpc}^{-3}$, our estimate translates to a rate $R \approx 3 \times 10^{-5}$ of TDEs yr$^{-1}$ per galaxy, which falls in between the estimate of Donley et al. (2002), $9.1 \times 10^{-6}$ yr$^{-1}$ per galaxy, and the one by Esquej et al. (2008), $2.3 \times 10^{-4}$ yr$^{-1}$ per galaxy. One can also try to predict this rate by combining equation (29) in Wang & Merritt (2004) for the expected dependence of the TDE rate on $M_{BH}$ and $\sigma$ with the well-known $M_{BH} \sim \sigma$ relation (e.g. Ferrarese & Ford 2005; Gültekin et al. 2009): $R \sim 3 \times 10^{-4}$ yr$^{-1}$ per galaxy. The fact that the rates derived from the cited observations are all well below this prediction could be due to a significant underestimate of the true intrinsic column density (as discussed by Sembay & West 1993; see the corresponding discussion in Donley et al. 2002 and Esquej et al. 2008), as well as due to uncertainties in the average properties of nuclear stellar clusters in nearby galaxies.

We should also mention attempts to systematically search for TDEs using long observations of predefined samples of galaxies, in particular clusters of galaxies (Cappelluti et al. 2006; Maksym, Ulmer, & Eracleous 2010; Maksym et al. 2013). Although such studies have found just a few TDE candidates, their results are also consistent with our estimated TDE rate of a few $10^{-5}$ yr$^{-1}$ per galaxy (Maksym et al. 2013).

5 CONCLUSIONS

We presented the results of a systematic search for tidal disruption events by looking for large (more than ten-fold) flux drops for sources from the ROSAT Bright Source Catalogue during serendipitous observations performed 10–20 years later by XMM-Newton. Besides a number of highly variable persistent X-ray sources (AGN, CVs and stars), we have found up to three sources that could be associated with TDEs in non-active galaxies. Specifically, three candidates are broadly consistent with the expectations of the TDE scenario: 1RXS J114727.1+494302, 1RXS J130547.2+641252 and IRXS J235424.5-102053, although the last object may be a spurious RASS source due to contamination by diffuse cluster emission. For 1RXS J112312.7+012858, a TDE association is also acceptable, but an AGN origin cannot be ruled out conclusively either. This implies the mean TDE rate $R \sim 3 \times 10^{-5}$ yr$^{-1}$ per galaxy within the surveyed volume, which is broadly consistent with previous estimates.

Due to the high sensitivity achieved in serendipitous observations by XMM-Newton, our sample of TDE candidates is significantly deeper than the previous ones by Donley et al. (2002) and Esquej et al. (2008). For this reason, it enables probing the TDE rate out to a larger redshift, up to $z \sim 0.18$ (making some reasonable assumptions about typical TDEs). The consistency of the TDE rate found in our work with the estimates based on the previous studies, which probed the more local Universe, tentatively suggests no significant evolution of the TDE rate between $z = 0$ and $z \sim 0.2$. The TDE rate derived here is also consistent with the rough estimate made for TDEs in clusters of galaxies (e.g. Maksym et al. 2013), thus indicating no dependency on the large-scale environment either.

We stress again that, in the absence of direct evidence, all of our suggested TDEs should still be regarded as TDE candidates. Further investigation of these sources is required (first of all, by means of optical spectroscopy) in order to confirm the absence of nuclear activity and to determine the distance to the host galaxy. Such information would help improve our estimate of the TDE rate. In addition, distance information would make it possible to estimate the peak luminosities of the TDEs and thus to constrain the masses of the SMBHs in the TDE host galaxies.

A by-product of our study is a small sample of high-amplitude AGN flares. More detailed studies of these sources could provide some insights into the mechanisms of AGN variability and consequently the physics of accretion onto SMBHs (like in the case of the Dynasty Foundation. The research made use of grant RFBR 13-02-01365.

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