Dissipative Properties of EHD Lubricant Film

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**Abstract:** For the case of the failure of the lubricant film at hydrodynamic lubrication a common thermodynamic theory of strength is considered. According to this theory the failure occurs when the internal energy density (potential and thermal components) in the volume of material reaches a constant for a given material. A special case of this theory is considered when only the density of heat (kinetic) component of internal energy is taken into account. Temperature condition determines the limit state for liquid lubricants - mineral oils. When analyzing the regularities of friction at EHD lubrication the state and properties of the oil film at the condition of irregular and hydrostatic compression. The original structural model of oil film at EHD lubrication in the form of the rotary oscillating cells with elastic interactions to each other is proposed. It is similar to the Rayleigh-Benard cells and corresponds to the cellular hypothesis of J. Gibbs for the case of equilibrium and reversible process. It is quite possible that the size of the cells have an order of about nano level. The oil film dissipates energy in the direction of relative motion of bodies. This oil film has the highest dissipative properties.

**Keywords:** friction, energy, dissipation, nanostructure, lubricant film

1. **Introduction**

We all know the friction is a global nature phenomenon of the energy transformation. Friction is subjected to energy balance equation and with thermodynamic point of view [1, 2] is the process of two interrelated, oppositely directed and concurrent trends operating in a strained contact. According to the energy balance scheme (figure 1) for plastic deformation and fracture [3-5] presented below, equations for friction work \( W_f \), frictional force \( F \) and friction coefficient \( \mu \) (without lubrication) has view:

\[
W_f = \Delta U_e + Q = \Delta U_{e_1} + \Delta U_{e_2} + \Delta U_{T_1} + \Delta U_{T_2} + \dot{Q}_1 + \dot{Q}_2, \tag{1}
\]

\[
W_f = \dot{U}_e + \dot{Q} = \dot{U}_{e_1} + \dot{U}_{e_2} + \dot{U}_{T_1} + \dot{U}_{T_2} + \dot{Q}_1 + \dot{Q}_2, \tag{2}
\]

\[
F_l = \frac{\Delta U_e}{l} + \frac{Q}{l} = \frac{\Delta U_{e_1} + \Delta U_{e_2}}{l} + \frac{Q_1 + Q_2}{l}, \tag{3}
\]

\[
F_v = \frac{\dot{U}_{e_1} + \dot{U}_{e_2}}{v} + \frac{\dot{Q}_1 + \dot{Q}_2}{v} = F_{\text{mechanical}} + F_{\text{molecular}}, \tag{4}
\]

\[
\mu = \frac{\Delta U_{e_1} + \Delta U_{e_2}}{Nl} + \frac{Q_1 + Q_2}{Nl} = \mu_{\text{adapt}} + \mu_{\text{dis}} = \mu_{\text{adapt}} + \mu_{T(\text{dis})} + \mu_{\dot{Q}(\text{dis})}. \tag{5}
\]
\[
\mu_v = \frac{\dot{U}_{e_1} + \dot{U}_{e_2}}{N_v} + \frac{\dot{Q}_1 + \dot{Q}_2}{N_v} = \mu_{\text{deformation}} + \mu_{\text{adhesion}},
\]

(6)

where \( \Delta U_e = V_f \Delta u_e \); \( Q = V_f q \); \( \dot{U}_e = V_f \dot{u}_e \); \( u_e = \text{da}_e/\text{dt} \); \( V_f \) - is the deformable (friction) volume; \( \mu \) - friction coefficient; \( \mu_{\text{adapt}} \) - adaptive friction coefficient; \( \mu_{T(\text{dis})} \) and \( \mu_{Q(\text{dis})} \) - static and dynamical components of dissipative friction coefficient; \( \Delta U_f \) - thermal component of internal energy; \( N \) - normal load; \( l \) - distance of friction. The latent energy density \( \Delta u_e \) is an integral parameter of tribostate and damageability (failure \( \Delta u_e^* \)).

Thus, viewed thermodynamically, the work done by friction forces \( W_f \) (the friction power \( \dot{W}_f \)), the friction force \( F \) and the friction coefficient \( \mu \) may be classified conventionally into two specific components with different kinetic behaviour [3-5].

The first component is associated with microscopic mechanisms of adaptive type and relates to the change of latent (potential) energy \( (\Delta u_{e_1}, \Delta u_{e_2}) \) of various elementary defects and damages that are generated and accumulate in the deformable volumes of materials friction pair (figure 2).

This energy is a unique and integral characteristic of the submicro–and microstructural transformations that occur in plastically strained materials [3-5]. This energy is a measure of strain hardening and damageability of materials.

**Figure 1.** Scheme of the energy balance for the plastic deformation of a solid body [3-5].

**Figure 2.** Schematic view of elementary friction’s contact [1, 2].

The second component is associated with microscopic mechanisms of dissipative type and relates to dynamic recovery processes in which latent energy and frictional heat are released \( (q_1, q_2) \). This energy originates in the motion and destruction of various elementary defects of opposite signs, the egress of these defects to the surface, the healing of reversible submicroscopic discontinuities, etc.
The ratios of the components $\Delta u_{c_1}$ and $\Delta u_{c_2}$ as well as $q_1, q_2$ of the balance vary over a wide range, depending on the physical, chemical, and structural properties of the materials that comprise the friction couple and the friction conditions. During this transformation the friction surfaces are heated and the heat of friction is transferred to oil film. In the course of time an oil film may fail.

2. Energy criterion of oil film failure

2.1. Thermodynamic theory of strength

Thermodynamic condition [3, 4] for local failure has view

$$u(\overline{r}_*, t_0) = u(\overline{r}_*, 0) + \int_0^t \dot{u}(\overline{r}_*, t) \, dt = u_* = \text{const}. \quad (7)$$

Here $u(\overline{r}_*, 0)$ - represents the density of the internal energy within the local macrovolume of the material in the initial (prior to deformation, $t=0$) state; $\dot{u}(\overline{r}_*, t)$ - represents the specific power of the internal energy sources in the local volume responsible for the failure; $\overline{r}_*$ - is the parameter which characterizes the coordinates $(x_*, y_*, z_*)$ of the local volume responsible for failure.

In accordance with the modern thermodynamic concepts relating to the strength and fracture of materials [3,4], having a fundamental nature, failure of a volume of a material (critical "defectness") occurs when the internal energy density $u$ (potential $u_c$ and thermal $u_T$ components) in this volume reaches the critical value $u_*$, constant for a given material. The criterion $u_*$ is a single-valued and integral characteristic of limiting damage ("defectness") of a material. Thermal failure of the material is the particular case of this theory when the change of the potential component of the internal energy is negligibly small and can be neglected.

The temperature, however, is a quantitative measure of macroscopic manifestation of change of the density of the thermal (kinetic) $u_T$ component of the internal energy of the material. We shall use these concepts to describe the case of oil film failure during friction with account for the kinetic characteristics of internal energy accumulation by liquid materials (oils).

2.2. Thermodynamic failure condition for solids

This condition [3,4] has the form

$$u = u_0 + \Delta u = u_* \quad (8)$$

Here

$$u_0 = u_{c0} + u_{T0}, \quad (9)$$

$$\Delta u = \Delta u_c + \Delta u_T. \quad (10)$$

$u$ - density of internal latent energy in volume of material; $u_{(0)}, \Delta u$ - internal energy of material in initial state and its change during friction; $u_{c(0)}, u_{T(0)}$ - potential and thermal components of internal energy of material in the initial state ($t=0$); $\Delta u_c, \Delta u_T$ - change of potential and thermal components of internal energy of material during friction.

2.3. Thermodynamic failure condition for liquid oil films

We can say the liquid materials cannot accumulate the energy of elementary defects as solids. Therefore, in the particular case for liquid oil films $u_{c0}$ and $\Delta u_c$ are equal to zero, then the failure condition (8) can be written in the form

$$u_T = u_{T0} + \Delta u_T = u_*^{T}. \quad (11)$$
Thus, equation (11) is the energetic condition of oil film failure in the contact [6]. In accordance with this equation the oil film fails if the density of the thermal (kinetic) energy $u_T$ in its volume reaches the critical value $u_T^*$. Upon reaching this value the oil loses its lubricating properties and friction transitions to the regime of friction without lubrication. As the condition of the absence of scuffing as a consequence of lubricant film failure we take the relation

$$u_T = u_{T0} + \Delta u_T < u_T^*.$$  

(12)

Dividing both sides of (12) by the oil density $\rho$ and the average oil heat capacity $\bar{c}_p$, we obtain

$$T = T_0 + \Delta T < T_{cr} = T_*.$$  

(13)

Here

$$u_{T0} = \rho \bar{c}_p T_0 = \int_0^{T_0} \rho c_p dT;$$

$$\Delta u_T = \rho \bar{c}_p \Delta T = \int_{T_0}^{T_*} \rho c_p dT;$$

$$u_T^* = \rho \bar{c}_p T_* = \int_0^{T_*} \rho c_p dT.$$  

(14)

$T_*$ - temperature of oil film failure; $T_0, \Delta T$ - surface temperature and temperature flash.

Equation (13) is the well-known H.Block relation [6] for seizure, where $T_{cr}$ is the critical temperature, upon reaching which the oil loses its lubricating properties and seizure of the surfaces takes place.

**Table 1.** Energetic Characteristics of Mineral Oils [7].

| Oil                        | $\rho$, kg/m$^3$ | $T_{cr}$, $^\circ$C [9] | $u_T^*$, MJ/m$^3$ |
|----------------------------|------------------|--------------------------|------------------|
| Instrument Vaseline, GOST 1805-51 (MVP) | 878     | 100                     | 617              |
| Transporter, GOST 982-50   | 880              | 180                     | 751,6            |
| Spindle, GOST 1642-50 (AU) | 891              | 165                     | 756              |
| Compressor, GOST 5546-54 (KhF-12) | -          | 160                     | 727,4            |
| Turbine, GOST 32-53 (ZOUT) | 895              | 120                     | 663,2            |
| Medicinal Vaseline, GOST 3164-52 | 875     | 20                      | 483,4            |
| Turbine, GOST 32-53 (ZOUT) | -                | 140                     | 693              |
| Industrial, GOST 5289-51(50)| 906              | 140                     | 705,5            |
| Diffusion, GOST 7904-56 (DI) | -                | 20                      | 491,7            |
| Autotractor, GOST 1862-42 (AK-10) | 920     | 150                     | 733,6            |
| Autotractor, GOST 1862-42 (AK-15) | 930     | 140                     | 724,2            |
| Naphthene-paraffin fraction of MS-20 oil (NPS MS-20) | -          | 155                     | 718,2            |
| Aviation, GOST 1013-49 (MZS) | 890              | 140                     | 693              |
| Aviation from Groznenskll crude, GOST 1012-29 | 905       | 210                     | 824,2            |
| Cylinder, TUM NP 233-47 (bright stock) | -            | 140                     | 693              |
| Ethylene glycol            | 1110            | 125                     | 833              |

Table 1 presents the values of $u_T^*$ for known values of $T_{cr}$ for mineral oils [8].
3. The calculation correlation with scuffing terms under lubricated friction

To establish the connection between the energetic characteristics of oil film failure in the contact of rubbing bodies and the friction process parameters we can use the relations presented in [9] for calculating the average temperature rise in the sliding contact friction zone. Since in the case of friction with lubrication the temperature of the thin lubricant film reflects the thermal state of the friction surfaces, then with accuracy adequate for practical purposes the values calculated for the friction surface temperature can be referred to the thermal state (temperature) of the oil film as well. Thus, to calculate the average temperature rise in the friction zone for sliding contact on an area in the form of a square with side \( 2l \) we take a relation [9], which at high sliding speeds has the form

\[
\Delta T = \frac{1.064\omega}{\lambda_1} \left( \frac{2a_1}{v} \right)^{1/2}.
\]  

In this case the relations (15) transform to the form

\[
\Delta u_T = \frac{1.064\omega_{lub}^*}{\lambda_1} \left( \frac{2a_1}{v} \right)^{1/2} ;
\]

\[
u_1^* = u_{T0} + \frac{1.064\omega_{lub}^*}{\lambda_1} \left( \frac{2a_1}{v} \right)^{1/2} ;
\]

\[
\omega_{lub}^* = \frac{(u_1^* - u_{T0})}{1.064\rho \bar{c}_p} \left( \frac{v}{2a_1} \right)^{1/2}.
\]

Here \( \omega_{lub} \) - specific friction power in friction with lubrication; \( \rho \) - density of oil; \( \bar{c}_p \) - average heat capacity of oil; \( l \) - length of contact spot; \( a_1, \lambda_1 \) - thermal and heat conductivity of bearing material; \( \lambda_2 \) - heat conductivity of counter-body material.

The condition (18) is the quantitative criteria characterizing the moment of transition from friction with lubrication to the friction regime without lubrication. Because of breakdown of the lubricant film, i.e., upon satisfaction of the condition \( \omega_{lub} = \omega_{lub}^* \) the oil loses its lubricity - the ability to intensely reduce the surface energy (passivate the surface) - and the friction regime without lubrication takes place.

Analysis of equation (18) shows that the scuffing criterion \( \omega_{lub}^* \) depends significantly on the critical value of the oil internal (thermal) energy density \( u_1^* \) and on the sliding speed (other conditions being the same). The larger the value of \( u_1^* \), the larger the specific friction power \( \omega_{lub}^* \) at which oil film failure occurs, which creates conditions for increase of the area of real contact of dry unlubricated surfaces. The influence of increase of the sliding speed \( v \) on increase of \( \omega_{lub}^* \) shows up as a consequence of increase of the volume (flow rate) of the lubricant pumped through the friction zone, which leads to improvement of heat transfer and as a result of this to reduction of the temperature in the friction zone and the parameter \( u_{T0} \). The thermo physical properties of the shaft and ring materials \( a_1, \lambda_1, \lambda_2 \) have an influence on the parameter \( \omega_{lub}^* \).

4. HD experiment and results

The experimental evaluation of the energetic condition of oil film failure during friction was performed on a full-scale stand with a pulsing load created by the hydraulic technique, which makes it possible to simulate the conditions of operation of the diesel engine connecting rod sliding bearings [7, 10]. The tests were made with the actual bimetallic sliding bearings (bushings) with a layer of antifriction material of the alloys AlSn20-1 and CuPb30 (\( D = 80 \) mm, \( H = 34 \) mm, \( h_{B} = 2.5 \) mm). Here \( H \) - bearing width; \( \varepsilon \) - diametral clearance; \( h_{B} \) - bearing thickness. The sliding bearings were tested in a pair with a steel
40NiCrMo6 (DIN) (induction hardened, HRC 50-55) shaft. Mark SAE40 diesel oil was used to lubricate the bearings. The experiments were made with shaft sliding velocity 4 m/sec. Special design of thermocouple has determined the real temperature into oil jet (stream) flowing out of connecting rod bearing (figure 3).

The limiting load capacity of the friction pair was determined on the basis of marked increase of the friction power $W_f$ and the temperature $T_{oil}$ of the oil leaving the connecting rod bearing (figure 3) in the given loading stage $p_{cr}$, which was taken as the limiting (critical) value. The results of the experiments are shown in Table 2.

![Figure 3. The disposition of thermopairs for measuring of temperature $T_{oil}$ of the oil leaving the connecting rod bearing.](image)

**Table 2.** Results of Experimental Studies of Load Capacity of AlSn20-1 and CuPb30 Alloys on Full-Scale Test Stand [10].

| Material     | AlSn20-1 | CuPb30 |
|--------------|----------|--------|
| $p_{cr}$, MPa| 48,0     | 35,0   |
| $W_f$, W     | 5157     | 9375   |
| $\mu$        | 0,011    | 0,027  |
| $T_{oil}$, °C| 120      | 115    |
| $A_{exp} \cdot 10^{-4}$, m² | 14,4     | 12,5   |
| $A_{calc} \cdot 10^{-4}$, m² | 13,3     | 11,9   |
| $\omega_{lub}^*$, MW/m² | 3,6      | 7,5    |

From the obtained experimental data ($\omega_{lub}^*, l, v$) we determined the oil film internal energy density change $\Delta u_f$ in the shaft-bushing contact at the moment of scuffing from formula (16); the friction surface temperature flash $\Delta T_f \equiv \Delta T_{oil}$:
\[ \Delta T_{\text{oil}} = \frac{\Delta u_T}{\rho e_p}; \]

the running value of the oil film internal (thermal) energy density

\[ u_T^* = u_T + \Delta u_T; \]

the critical value of the oil film internal (thermal) energy density and the friction surface temperature at the moment of scuffing

\[ T_{cr} = \frac{u_T^*}{\rho e_p}. \]

Here \( T_{oil}, \Delta T_{oil} \) - running value of oil temperature and critical value of oil temperature change in volume of oil film.

In the calculations it was assumed that the temperature \( T_{oil} \) of the oil leaving the connecting rod bearing reflects the average bearing friction surface temperature.

The results of calculation of the oil film energetic characteristics and the friction surface temperature are shown in table 3. For the given diesel oil grad (table 3) the critical values \( u_T^* \) and \( T_{cr} \) are, respectively, 733 MJ/m\(^3\) and 150\(^\circ\)C.

**Table 3.** Results of Calculation of Oil Film Energetic Characteristics Corresponding to Moment of Scuffing [10].

| Material      | AlSn20-1 | CuPb30 |
|---------------|----------|--------|
| \( u_T \)     | 682      | 673    |
| \( \Delta u_T^* \) | 43       | 54     |
| \( u_T^* \)   | 725      | 727    |
| \( T_{oil} \) | 120      | 115    |
| \( \Delta T_{oil} \) | 25,5   | 31     |
| \( T_{cr} \)  | 145,5    | 146    |

Thus the experiment results confirm reliably the assumed energetic scuffing condition (11). The critical magnitude of the oil film internal (thermal) energy density for a given inactive mineral oil grade is independent of the bearing material grade and the friction process conditions and is a physical constant of the tribosystem scuffing process.

The obtained results agree well with the fundamental concepts of the ergodynamic theory of strength [3-5], in accordance with which the integral measure of material susceptible to damage is the critical (constant for a given material) value of the internal energy density, and also with the H.Block hypothesis [6] on constancy of the overall seizure temperature for a given mineral oil grade.

5. EHD lubrication hypothesis

It is known that the EHD lubricant type has significantly higher performance (capacity for work), than HD lubricant type. Modern tribology [11] high performances lubricant when EHD explains with the effect of increasing the viscosity of the lubricant when it is compacted. It is known that for most mineral oils use high pressures can lead to a huge increase in viscosity. The oil between the surfaces at high pressures in the contact region behaves virtually like a solid separating layer. This process explains why many mechanisms in practice operate under much severer conditions than the classical theory would allow.
Naturally the question arises about the viability of EHD lubrication in its temperature aspect. Is it possible to apply temperature failure criterion of an oil film that successfully explains the loss of efficiency at HD lubricant?

We can examine the Hertzian scheme of contact under friction (figure 4).

In proximity of Hertzian contact with size of a material point A due to elasticity (Poisson effect) two equilibrium (elastic) hemisphere (roughness) - on roller and on the surface are formed. If there is a thin film of lubricant between the surfaces this two roughnesses operate as elements of a gear pump. Roughness of roller captures lubricant, which is getting into the zone of Hertzian contact of material point A and is hydrostatically compressed and then pressed out into disclose contact on surface roughness. Here, we may consider these three stages of formation of the equilibrium contact roughnesses at EHD lubrication on an example of mechanical model of EHD lubricant (figure 5, it is turned by 180° to figure 4).

![Figure 4](image1.png)

**Figure 4.** Model of elastic perfectly smooth roller on elastic and smooth surface.

In the first stage (a) an oil film is captured with elastic roughness into the cavity under the roughness. In the second stage (b) an oil film is hydrostatically compressed in this cavity. At the third stage (c) a narrow gap begins to form between these two roughnesses and compressed lubricant flows at high speed from the field of hydrostatic compression, creating a peak of pressure and then dropping the stresses in oil film up to zero.

![Figure 5](image2.png)

**Figure 5.** Three stages of forming contact with EHD lubrication.

Proposed model for Hertzian lubricant contact is fully corresponds with the known curve of pressure distribution in contact EHD (figure 6).
Left and central parts of pressure distribution in elasto-hydrodynamical film of lubricant correspond to the areas of capture and compression. To the right of pressure peak part corresponds to condition of high-speed flow of compressed lubricant through a narrow gap on the third stage of forming contact of EHD lubrication.

The question about the nature and state of the oil film in EHD lubrication is directly linked to the question of temperature fracture criteria for lubricant. The logic is here quite simple. Why does under EHD lubrication where the loads exceed many times the load of HD lubricant, lubricating film remains healthy?

When analyzing the friction regularities at EHD lubrication it is necessary to pay attention to the two independent and possible effects. The first one is regularity of plastic deformation (states and properties) at Hertzian contact of solids. The second one is the state and properties of the oil film in the condition of irregular and hydrostatic compression.

Let us consider the first. It is deformation of the Hertzian contact. Hertzian contact of friction should be considered as a contact in the most fully developed evolution of plastic deformation. The degree of structural adaptations of such contact can be maximum [12]. In essence, we can have a perfect elasto-visco-plastic deformation. Such contact has the maximum capacity for work and temperature effect on the surface is minimal. Almost totally equilibrium (rotary elastically, interactive) nano-crystal structures of contact (mechanical quanta (figure 7) [13-16]) have thermal effect in the form of temperature growth on the temperature level is close to absolute zero. Such microflash is unable to heat and destroy the oil film. Irrelevance of lubricants is very often supposed for such contacts [11]. But on the other hand the problem of contact temperature cycling and the possibility of gradual heat accumulation require a lubricant which will remove the heat.

Figure 6. Model of Ertel-Grubin EHD film: $p_H$ - Hertzian pressure; $p(x)$ - current pressure layer; $h(x)$ - the current thickness of the lubricating layer; $a$ - half width of contact area.

Figure 7. Model of elementary dissipative nanostructure of friction (8103 atomic cubical cells) [13].

Thus, despite of the fact that Hertzian contact works at stresses at the threshold of the strength limit, within the contact between structural elements (mechanical quanta, as subtribosystems) interaction is realized at the level of the perfect elasticity. Elastic mutual rotations of nanostructures define the essence for models as abnormally low friction and wearlessness [14]. Such state of tribocontact we may examine as state of selforganization, nanoquantum solid lubricant [15].

Now, as to the second - it is about state and properties of lubricants at EHD lubrication. Quasi solid state of the oil film in full irregular hydrostatic compression can give unique lubrication properties.

First, let’s take into account the classic force-separation curve of interatomic interaction (figure 8). Hydrostatic compression of lubrication film makes the destruction process more difficult (lubricant state under point P). It is so, because stretching (attractive force, see point Q) is necessary for destruction (breaking of atomic bond).

Secondly, quasisolid state of oil film may analyze from the perspective of solids friction.
Therefore, the regularities of shear deformation of lubricant (oil) may be similar with regularities that take place under friction of solids, namely, one may remember the structural-energy diagram (figure 9) for evolution of the rubbing surfaces [1, 2, 12, 16].

**Figure 8.** Sketch, showing the way in which the force between two atoms varies with the distance between them.

**Figure 9.** Structural-energy diagram for evolution of rubbing surfaces [1, 2, 12-16].

Here may well work mechanism for dynamic dissipation of energy with energy balance component $\bar{q}^*$ (figure 1), i.e. EHD lubricant can have properties of dissipative structures of friction such by friction of solids. These dissipative structures can be formed at the gap mouth of the expiration of oil where the lubrication hydrostatic compression has irregularity at the expense of stretching (figure 10). Between the second and third stages of the process occurs a momentary restructuring lubricant from the most compact to a new nanostructured one. A phase transition takes place. A volume front of thermal (temperature) blow is momentarily formed at that. Actually this blow is a cause the formation of maximum internal pressure peak that precedes the development and process of the third phase –
laminar flow of nanostructured lubricant in the narrow clearance. At that a resistance of lubricant flow drops to zero. Internal stresses drop to zero that the chart of Ertel-Grubin is demonstrated. This condition is a perfect, elastic, inner balance between structural (nano) elements of lubricant.

**Figure 10.** Dissipative oil structures zone under EHD lubrication.

It momentary and volume front of temperature blow can’t warm up the lubricant film. It is only mechanism of energy transformation for transition the lubricant film state from state with latent elastic energy $\Delta u_e^{*}$ of hydrostatic compression to state with latent elastic energy $\tilde{q}_e^{*}$ of nanostructured flow. Practically all accumulated energy $\Delta u_e^{*}$ of the first stage is converted to energy $\tilde{q}_e^{*}$ of the third stage owing to the structural transformation of the oil film. Correspondingly, high density of energy flux (friction power) at EHD lubrication stimulates the formation of similar structures of quick heat dissipation. Here it is quite possible to propose structural model of oil film (figure 11) with EHD lubrication in the form of rotary oscillating cells with elastic interactions with each other.

**Figure 11.** Cellular (nano-structured) model of EHD lubricant into a narrow gap between two surfaces.

**Figure 12.** Rayleigh-Benard cells.

This is similar to the Rayleigh-Benard cells (figure 12) in the case of a powerful stationary heating of liquids. For static heating condition the Rayleigh-Benard cells pass the heat in perpendicular direction to friction surfaces. It is possible that the size of these cells has an order is about nano-scale level such as fullerene' form or mechanical quantum (figure 7) [12-16]. It corresponds to the cellular hypothesis.
of J. Gibbs for case of equilibrium and reversible process. Accordingly, this oil film has maximum dissipative properties.

At EHD lubrication this high structured oil film dissipates energy in the direction of relative motion of bodies (figure 11).

Here all accumulated energy $\Delta u_*$ in compressed zone (the work of external forces) is almost entirely scattered on the dissipative structures (subtribosystems) of oil film.

6. Conclusions

Thus, neither the first nor the second effects Hertzian contact do not give the temperature effect of friction. When EHD lubrication there is we have no temperature effect at destruction of lubricant.

The criterion of the ultimate state of oil film at EHD is a critical density $\bar{q}^*$ of dissipation energy, which is equal to the critical density of the internal thermal energy $u_T^*$ of lubricant. The concept of temperature for such film has almost no meaning. Here you should write ratio

$$ \Delta u_T = u_T^* - u_{T0} = 0. \quad (23) $$

This is associated with a particular (unique) structural state of oil film. The presence of dissipative component $\bar{q}^*$ presupposes the equality condition of zero temperature flashes $\Delta T = 0$ and consequently $\Delta u_T = 0$ too.

Therefore, the critical state of the film cannot be achieved until the dissipative structures lubricating film operate. Loadability of EHD lubrication film is very high because of the elastic and reversible energy dissipation.

As a result, we can write the following - the equations for critical power of friction for HD and EHD lubricants:

$$ W_{HD} = \dot{V} \Delta u_T^* = \dot{V} \Delta u_T^* \rho (T^* - T_0), \quad (24) $$

$$ W_{EHD} = \dot{V} \cdot \bar{q}^*. \quad (25) $$

Here $\dot{V}$ is the amount of lubricant in contact at a time.

This result is fully in line with the practice. EHD lubrication capacity for work should be significantly higher than HD lubrication

$$ W_{EHD} \gg W_{HD}, \quad (26) $$

Because

$$ \bar{q}^* \gg \Delta u_T^*. \quad (27) $$

If you write down the numerical values $\Delta u_T^*$ and $\bar{q}^*$ for engine oil, for example, M14V (Table 3) then we get the following experimental results for AlSn-20 and CuPb30 (in brackets):

$$ \Delta u_T^* = 43(54) \text{MJ/m}^3, $$

and

$$ \bar{q}^* = u_T^* = 725(727) \text{MJ/m}^3. $$

In the most general case, the thermal model of the ultimate state of oil film with HD and EHD lubricants have the following forms:

HD lubricant - $\dot{W}_{HD} = \dot{V} q^*$, here, $q^* = \Delta u_T^*$, and $\bar{q}^* = 0$;

EHD lubricant - $\dot{W}_{EHD} = \dot{V} q^*$, here, $q^* = \bar{q}^* = u_T^*$, and $\Delta u_T^* = 0$.

Excess of specific friction power above critical level would lead to turbulence of the flow, laminar flow destruction of dissipative structures and consequently to loss of load capacity of EHD. It will be a critical state.
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