Localization of inclusions in multiple prompt gamma ray analysis: a feasibility study

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Abstract. We investigate the feasibility of using low energy gamma rays from neutron capture to localize slabs inside samples. A new system based on two gamma ray detectors with 2D collimators to be tested at the INES beamline at the pulsed neutron source ISIS (Oxford, UK) is described. The system provides a localization of slabs inside extended samples by using gamma ray self-absorption. Geant4 Monte Carlo simulations of the beamline were carried out to model gamma spectra from test samples.

1. Introduction

Neutron capture prompt gamma activation analysis (PGAA) is a non-destructive analytical technique that provides information on the isotopic composition in samples. It is based on the measurements of gamma rays following neutron capture. The pattern of gamma rays, ranging from hundred keV to 12 MeV, is unique for each isotope [1]. The PGAA method has been applied in materials science, chemistry, geology, mining, archaeology, environment, food analysis, and medicine [2]. A current limitation of the technique is gamma ray self-absorption, causing attenuation of gamma rays within the sample itself. That may lead to significant errors [1]. In this paper, we exploit the gamma ray self-absorption and investigate the feasibility of using low energy gamma rays (< 650 keV) to localize metal slabs inside extended objects. Depending on the position of the inclusion, the gamma rays will be attenuated differently. Using two back-to-back detectors on opposite sides of the sample we can measure the ratio of the gamma ray intensities and calculate the approximate position of the slab. Spatial distribution of elements has recently been measured by scanning a sample with a collimated neutron beam and by using a collimated gamma detector [3]. Drawback of the technique is the high collimation of source and detector needed and the long beamtime required. The method described in this paper will make use of the full neutron beam area and allow fast localization of the slabs along the direction where the gamma ray detectors are placed. Information on the position of the inclusion could be used to correct for the self-absorption and obtain more precise quantitative analysis. This study has been performed with Monte Carlo simulations based on the GEANT4 code [4, 5].
2. Materials and Methods

2.1. Theory of measurements

The attenuation of gamma rays within the sample increases with thickness, density, and atomic number of attenuating medium, and with the decreasing of the gamma ray energy [1]. For gamma rays penetrating a uniform layer of material in the normal direction, the transmitted flux is:

\[ I = I_0 \exp(-\mu x) \] (1)

where \( I_0 \) is the incident neutron flux, \( \mu \) is the attenuation coefficient, and \( x \) is the layer thickness. In a neutron capture event the gamma rays are emitted isotropically. In order to ensure that the gamma ray detectors measure intensities of gamma rays reaching the detectors after traveling along a straight-line path from the gamma ray source, i.e. the inclusion, we employed 2D collimators. The collimators were located immediately in front of the detectors. The ratio of counts on the gamma ray detectors A and B can therefore be expressed as:

\[
\frac{C_A}{C_B}(E) = \frac{I_0(E) \times \epsilon_A(E) \times \exp\left(-\sum_{i=1}^{N} \mu_i(E) \times x_i^A\right)}{I_0(E) \times \epsilon_B(E) \times \exp\left(-\sum_{i=1}^{N} \mu_i(E) \times x_i^B\right)}
= \frac{\epsilon_A}{\epsilon_B}(E) \exp \left[-\sum_{i=1}^{N} \mu_i(E) \times (x_i^A - x_i^B)\right]
\] (2)

where \( I_0(E) \) is the number of the gamma rays emitted with energy \( E \), \( \epsilon_A(E) \) is the efficiency of the detector A, \( \epsilon_B(E) \) is the efficiency of the detector B, \( \mu_i \) is the attenuation coefficient of the material \( i \), \( x_i^A \) is the distance traveled by the gamma ray in the material \( i \) in the direction of the detector A, and \( x_i^B \) in the direction of the detector B.
For samples in which the matrix is uniform and the slabs are small enough to ignore absorption within the slabs themselves, we can calculate the position of the slabs, knowing that \( x_A^A + x_B^B = D \), as:

\[
x^A = \frac{1}{2} \times \left[ D - \frac{1}{\mu_M(E)} \ln \left( \frac{C_A(E) \times \epsilon_B(E)}{C_B(E) \times \epsilon_A(E)} \right) \right]
\]  

(3)

where \( \mu_M \) is the attenuation coefficient of the matrix, and \( x_A^A \) is the distance of the center of mass of the slab from the sample edge along the \( x \) axis toward the detector A, \( D \) is the sample thickness in the direction \( x \). The position of the slab will be:

\[
x_{slab}^A = \frac{D}{2} - x^A = \frac{1}{2 \times \mu_M(E)} \ln \left( \frac{C_A(E) \times \epsilon_B(E)}{C_B(E) \times \epsilon_A(E)} \right)
\]  

(4)

The values of \( x_{slab} \) are independent of the gamma ray energy. For that reason, to decrease uncertainties we averaged on all the values of \( x_{slab} \).

This procedure was applied to several samples of Fe and Al with slabs of thickness 0.1 cm at different positions.

2.2. Monte Carlo simulation

The proposed multi prompt gamma ray analysis system consists of two back-to-back gamma ray detectors and two 2D collimators located immediately in front of the detectors. Detectors diameter and length were 6.4 cm and 9 cm, respectively. The gamma ray detectors were placed perpendicular to the primary beam at opposite sides of the sample at a distance of 15 cm. The 2D lead collimators were 2.2 cm thick with 0.1 cm septa and 0.4 cm periodicity. Figure 1 shows the simulated geometry. The simulation was carried out for a water moderator at 293 K at a spallation source. The neutron energy spectrum is shown in figure 2. To speed up the simulation we selected neutron energies up to epithermal range (< 1.2 keV). The neutron beam was simulated as a parallel beam of size 2 x 2 cm².

![Figure 2. Neutron energy spectrum simulated with the Monte Carlo radiation transport code MCNPX (private communication with S. Ansell [6, 7]).](image-url)

Samples with one and two slabs were simulated to show the capability of the system to localize slabs inside a matrix representing a model of a realistic sample. The matrix size was 2 x 2 x 2 cm³. The matrix material was Fe or Al. The slab size was 0.1 x 2 x 2 cm³. The slab material was Cu and/or Fe. Neutron tracks were followed through the simulated geometry until they were absorbed or exited the world volume. Inelastic and elastic scattering were simulated with the high precision neutron model of Geant4 [8]. The neutron capture was modeled by randomly...
Figure 3. Simulated energy spectrum measured by detector A (left) and B (right) from a Fe sample with a Cu slab in the energy range [0, 800] keV.

Sampling gamma ray energies from the nuclear decay schemes (Evaluated Nuclear Structure Data File [9]). The radioactive decay of unstable nuclei produced by neutron capture was simulated with the Geant4 radioactive model [8]. Electromagnetic interactions were simulated with the Geant4 Livermore model [8]. Energy and measurement time were recorded for every gamma ray detected. To evaluate the feasibility of the method without loss of generality we considered $\epsilon_A(E) = \epsilon_B(E) = 1$ in equation 4. The number of simulated events was $10^9$.

3. Results

Figure 3 shows a simulated energy spectrum measured by detector A and B from a Fe sample with a Cu slab. Gamma ray lines from neutron capture in Fe and Cu are clearly visible in the energy range [0, 800] keV.

Figure 4 shows the ratio of intensities of Cu lines as a function of gamma ray energies for slabs in a Fe matrix placed at offset -0.2 cm and 0.2 cm along the x axis. The ratio $C_A/C_B$ is well separated for the two offsets for all the gamma ray lines.

Table 1 lists the reconstructed position of the slabs calculated from equation 4 for different samples. The error on the calculated Cu slab position is the standard deviation of the $x^{slab}$ position calculated for each peak. The Fe element shows only a gamma-ray peak, at 352 keV, in the selected energy range ($< 650$ keV) in accordance with the prompt gamma-ray spectrum catalog [1] in which only the 352 keV peak has partial gamma ray production cross section above 0.1%. The difference between the actual and reconstructed position of the slabs is $\leq 0.1$ cm for all the studied configurations. Samples containing less abundant and/or weaker gamma-emitters could be investigated using the neutron time of flight (TOF) information and the resonance structure of neutron capture cross section. Gamma ray lines of less abundant or weaker gamma-emitter isotopes could be discerned by selecting gamma ray spectra at specific TOF windows. In future, simulations accounting for detector effects and gamma-ray background will be performed in order to select the most suitable detector and to study limits and geometrical uncertainties of the technique for different matrix/slab materials and thicknesses. Measurements at the pulsed neutron source ISIS (Oxford, UK) are foreseen.

4. Conclusions

The feasibility of using low energy gamma rays from neutron capture to localize metallic slabs inside a sample has been investigated with Monte Carlo simulations. A detection system consisting of two back-to-back gamma ray detectors with two 2D collimators has been proposed.
The position of the slab has been correctly reconstructed using the ratio of the intensity of two gamma ray detectors placed on opposite sides of the sample. Simulations taking into account detector effects and studying the limitations of the techniques are ongoing. Information on inclusion position could be used to correct for the gamma ray self-absorption and to obtain more precise quantitative analysis. Possible applications of the technique are in metallography, geology, and archeology.

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References
[1] Molnar G L 2004 Handbook of Prompt Gamma Activation Analysis with neutron beams (Kluwer Academic Publishers) ISBN 1-4020-1304-3
[2] Perry L A, Firestone R B, Molnar G L, Revay Z, Kasztovszky Z, Gatti R C and Wilde P 2002 J. Anal. At. Spectrom. 17 32–37
[3] Belgya T, Szentmiklosi L, Kis Z, Kasztovszky Z, Schulze R, Kudejova P, Materna T, Canella L and Jolie J 2009 Proceedings of the 14th International Conference on Cultural Heritage and New Technologies
[4] Agostinelli S et al. 2003 Nucl. Instr. and Meth. A 506 250 – 303
[5] Allison J et al. 2006 Nuclear Science, IEEE Transactions on 53 270–278
[6] Bedogni R, Esposito A, Andreani C, Senesi R, Pascale M P D, Picozza P, Pietropaolo A, Gorini G, Frost C D and Ansell S 2009 Nucl. Instr. and Meth. A 612 143 – 148

[7] Rebai M, Gorini G, Perelli-Cippo E, Pietropaolo A, Tardocchi M, Fazzi A, Milani E, Verona-Rinati G, Andreani C, Senesi R, Frost C D, Schoonevled E M, Rhodes N J, Bedogni R and Esposito A 2010 XIX ICANS Conference Proceeding

[8] Geant4 collaboration 2011 Physics Reference Manual

[9] Tuli J K 2001 Evaluated Nuclear Structure Data File (Upton NY 11973-5000 USA)