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Lifetime of sessile saliva droplets in the context of SARS-CoV-2

Saravanan Balusamy\textsuperscript{a}, Sayak Banerjee\textsuperscript{a}, Kirti Chandra Sahu\textsuperscript{b,}\textsuperscript{*}

\textsuperscript{a} Department of Mechanical and Aerospace Engineering, Indian Institute of Technology Hyderabad, Sangareddy 502 285, Telangana, India
\textsuperscript{b} Department of Chemical Engineering, Indian Institute of Technology Hyderabad, Sangareddy 502 285, Telangana, India

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\textbf{ABSTRACT}

Spreading of respiratory diseases, such as COVID-19, from contaminated surfaces is dependent on the drying time of the deposited droplets containing the virus. The evaporation rate depends on environmental conditions, such as ambient temperature and relative humidity and physical properties (e.g., droplet volume, contact angle and composition). The respiratory droplets contain salt (NaCl), protein (mucin), and surfactant (dipalmitoylphosphatidylcholine) in addition to water, which are expected to influence the evaporation in a big way. A diffusion-based theoretical model for estimating the drying time is developed which takes into account the dynamic contact angle of saliva droplets laden with salt and insoluble surfactants. The effect of the initial volume, contact angle, salinity, surfactant concentration, ambient temperature and relative humidity on the drying time of droplets is investigated.

1. Introduction

SARS-CoV-2 (also known as coronavirus) has spread to more than 200 countries and infected more than 11 million people worldwide until the first week of July 2020. Respiratory infections such as SARS-CoV-2 virus spreads primarily by respiratory droplets (size $\leq$10\,\mu m) of saliva or discharge from the nose of an infected person during coughing/sneezing \cite{1,2}. The study of fluid dynamics of respiratory droplets come into contact with a surface depending on the surface properties and local ambient conditions. Bhardwaj and Agrawal \cite{4} recently used a diffusion-based model to study pinned droplet evaporation with a constant contact angle assuming that SARS-CoV-2 droplet is pure water. They also remarked that the quasi-steady assumption used in their study is valid when the heat equilibrium time is much smaller than the drying time of the droplet \cite{7}. Respiratory droplets, however, also contain salt (NaCl), protein (mucin), and surfactant (dipalmitoylphosphatidylcholine), which are supposed to influence the drying time. In particular, the concentration of NaCl, mucin and surfactant per litre of saliva solution are 9 g/l, 3 g/l and 0.5 g/l, respectively \cite{8}. Thus, in this study, we have extended the study of Bhardwaj and Agrawal \cite{4} to investigate the lifetime of sessile droplets laden with salinity and other insoluble ingredients (surfactant) on various surfaces under different atmospheric conditions, such as relative humidity (RH) and temperature (T), which has been poorly understood until now.

In the context of normal droplets, the evaporation has been investigated theoretically and experimentally by many researchers (see e.g. Refs. \cite{4,9-11}). Gurrala et al. \cite{9} investigated the evaporation of sessile droplets at different temperatures. They compared the results obtained from the theoretical model that includes the diffusion, free convection and passive transport of air with their experiments for different compositions and surface temperatures. Fukatani et al. \cite{10} experimentally studied the effect of the relative humidity and ambient temperature on the early stages of the evaporation of organic solvents. Karapetsas et al. \cite{12} investigated the evaporation of sessile droplets laden with particles and insoluble surfactants using lubrication theory. They found that the lifetime of the droplet is significantly influenced by the presence of the surfactant via the spreading of the droplet and Marangoni convection. Souli et al. \cite{13} compared the initial evaporation rate (linear stage) of droplets of aqueous saline (NaCl) solutions with that obtained from their experiments, and found that the evaporation rates for high salt concentrations are significantly affected by the presence of the surfactant.

\textsuperscript{*} Corresponding author.
\textit{E-mail address:} ksahu@che.iith.ac.in (K.C. Sahu).

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concentrations and small contact angles were significantly lower than those predicted from the well-accepted diffusion-controlled evaporation scenario for sessile droplets. They also observed that the droplet remains pinned when the concentration of the salt is \( \geq 10^{-6} \) mole.

2. Mathematical modelling

For small droplets (volume, \( V \leq 10^{-3} \) \( \mu \)l), the surface tension force predominates over the gravitational force and a sessile droplet obtains a symmetrical spherical cap shape [14]. For such a droplet, if \( A(t) = \theta(t) \cdot \pi/180 \) is the contact angle in radians and \( R \) is the wetting radius of the droplet, then the time-varying droplet volume can be expressed as

\[
V(t) = \frac{\pi R^3}{3} \left(1 - \cos A(t)\right) \left(2 + \cos A(t)\right) \theta(t)
\]

We make the following assumptions for the deduction of the rate of evaporation of droplets laden with salinity and insoluble surfactants. (i) The surface is at the same temperature as the surrounding air, (ii) the temperature remains constant during the evaporation process, (iii) the droplet remains internally well-mixed during the evaporation process and (iv) the Stefan flow and the effect of solutal Marangoni convection inside the droplet are neglected. These assumptions are valid for the evaporation of small droplets. Under these conditions, one can ignore the convective mass transfer and assume that the droplet evaporates at the vapour-liquid interface only through the processes driven by diffusion. Therefore, the well studied expression [15] of the evaporation rate, \( \dot{m} \) (kg/s) of spherical cap-shaped droplets through the diffusion process are adopted in this work, which is given by

\[
\dot{m} = -\pi R^2 \left(1 - \cos A(t)\right) \left(2 + \cos A(t)\right) \theta(t) \dot{f}(A) \alpha
\]

where \( \dot{f}(A) \) is the relative humidity of air (%) and \( f(A) = 1.3 + 0.27A^2 \) for \( \theta \leq 90^\circ \) [15]. The diffusion coefficient, \( \mathcal{D}(T) \) of water vapour (m²/s) at the given temperature, \( T \) (°C) is given by [4].

\[
\mathcal{D}(T) = 2.5 \times 10^{-5} \exp\left(\frac{-684.15}{T + 273.15}\right)
\]

For water-salt-surfactant solutions, the saturation vapour density, \( C_{\text{sat}} \) of water (kg/m³) over the solution interface can be obtained from the Raoult’s law,

\[
C_{\text{sat}} = X_w C_{\text{sat}}
\]

where \( X_w \) is the mole fraction of water in the solution, and \( C_{\text{sat}} \) is the saturation vapour density of pure water at the given temperature, which is given by [4].

\[
C_{\text{sat}}(T) = 9.99 \times 10^{-7}T^2 - 6.94 \times 10^{-3}T + 3.20 \times 10^{-2} - 2.87 \times 10^{-2}
\]

In Eq. (2), \( \alpha \) is the accommodation coefficient for evaporation of the droplet laden with insoluble surfactants [12], which is given by

\[
\alpha = \frac{1}{1 + \Psi TT_w}
\]

where, the surfactant parameter, \( \Psi \geq 0 \), and

\[
\Gamma = \frac{n_{\text{sur}}}{n_s(t) + \bar{n}_s + n_{\text{sur}}}, \quad \Gamma_w = \frac{n_w}{n_s(t) + \bar{n}_s + n_{\text{sur}}}
\]

Here, \( \Gamma \) is the concentration of the surfactant and \( \Gamma_w \) is the maximum possible surfactant concentration which is assumed to be the concentration corresponding to a fully dried drop. Also, \( n_{\text{sur}} \) is the number of moles of surfactant, \( \bar{n}_s \) denotes the number of moles of solute ion particles and \( n_{\text{sur}}(t) \) is the number of moles of water in the droplet at the given time. The number of moles of surfactant is calculated using initial value of \( \Gamma = 1.22 \times 10^{-5} \) [8]. The mole fraction of the water in the solution decreases as the evaporation process progresses over time. This decrease in the mole fraction of water in the evaporating droplet increases the surfactant and salt concentrations which, along with a decrease in the contact angle, slows down the rate of evaporation as time progresses. At any given time, the mole fraction of liquid water that remains in the solution can be calculated as

\[
X_w(t) = \frac{n_{\text{w},t}}{n_s(t) + \bar{n}_s + n_{\text{sur}}}
\]

Note that for electrolyte solutes that dissociate into ions, the sum of the number of moles of the positive and negative ions formed by the solute is equal to \( n_{\text{salt}} \). The molality (M) is defined the ratio between the number of moles of solute (\( n_{\text{salt}} \)) and mass of water (\( m_w \) in kg) in the droplet solution. For a salt, like NaCl, \( 
\bar{n}_s = 2n_{\text{salt}} \). Since salt and surfactants are non-volatile, the mass evaporation rate is entirely due to the loss of water. We neglect any loss of these ingredients due to crystallisation from the solvent as the initial droplet is very dilute and no deposition is expected till the very end of the droplet’s lifetime. Thus the rate of change of the number of moles of water from the droplet solution can be expressed as

\[
\frac{dn_{\text{w}}}{dt} = \frac{\dot{m}}{\mu_w}
\]

where, \( \mu_w \) is the molecular weight of water. The rate of change in the droplet volume is given by

\[
\frac{dV}{dt} = \frac{\dot{m}}{\rho(t)}
\]

where \( \rho(t) \) is the density of the solution at a given time. The fluid density is a function of the temperature and the salt and surfactant concentrations in the solution and its value is obtained from polynomial fits of appropriate thermodynamic data charts [16].

Fig. 1. Temporal variations of (a) the contact angle of the droplet normalised with the initial contact angle, \( \theta/\theta_0 \), and (b) the normalised instantaneous mass evaporation rate of the droplet with its initial mass evaporation rate, \( m/m_0 \) for different values of \( \Psi \). The rest of the parameters are \( V_0 = 10 \) nl, \( M = 0.154 \) mol/kg, \( T = 30^\circ\text{C}, \text{RH} = 50\% \) and \( \theta_0 = 50^\circ\). The inset in panel (b) shows a magnified view during the crossover. The dashed horizontal line represents the constant contact angle case [4].
3. Results and discussion

The conditions used in the present study are summarised below. The initial contact angle ($\theta_0$) of the droplet is varied from 10° to 90° to represent different surfaces used in our everyday activities. Note that the contact angle of a droplet on glass, stainless steel, cotton, wood, smartphone screen and coated paper are about 10° [17], 30° [18], 50° [19], 65° [20], 85° [21], and 90 – 110° [22], respectively. The relative humidity (RH) and temperature ($T$) are varied from 10% - 90% and 20°C - 50°C, respectively. Such ambient conditions are observed in various places worldwide. The initial volume of the droplet ($V_0$) is varied from 1 nL - 10 nL. In addition, the molality ($M$) of the saliva is fixed at 0.154 mol/kg [8]. For a typical value of $\Gamma/\Gamma_\infty$, the value of $\Psi$ varies from 0 to 50 for different realistic values of $\alpha$ [23].

In the present study, given the initial mole numbers of water, surfactant and solute ions in the droplet, an iterative numerical code is developed which calculates the mass loss rate while dynamically correcting the changes induced by the decrease in the contact angle and the increase in the salt and surfactant concentrations in the liquid during the evaporation process.

We begin the presentation of our results by examining the effect of the surfactants on evaporation by exploring the influence of the surfactant parameter $\Psi$. Fig. 1(a) and (b) show the temporal evolutions of the normalised contact angle ($\theta/\theta_0$) and the normalised instantaneous mass evaporation rate ($\dot{m}/\dot{m}_0$) of the droplet for different values of $\Psi$, respectively. The rest of the parameters are $V_0 = 10$ nL, $M = 0.154$ mol/kg, $T = 30$°C, $\Psi = 10$ and $\theta_0 = 50°$.

Fig. 2. Temporal variations of (a) the contact angle of the droplet normalised with the initial contact angle, $\theta/\theta_0$, (b) the normalised instantaneous mass evaporation rate of the droplet with its initial mass evaporation rate, $\dot{m}/\dot{m}_0$ for different values of RH (%). The rest of the parameters are $V_0 = 10$ nL, $M = 0.154$ mol/kg, $T = 30$°C, $\Psi = 10$ and $\theta_0 = 50°$.

Fig. 3. Regime maps showing the droplet’s drying time in the $T$-RH space. (a) $\theta_0 = 10°$, (b) $\theta_0 = 90°$. The colorbar shows the drying time in second in the logarithmic scale. The rest of the parameters are $V_0 = 10$ nL, $M = 0.154$ mol/kg and $\Psi = 20$. 

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of water nonlinearly. Both the Fig. 1 (a) and (b) shows a highly nonlinear information shows that the corresponding decrease in the mole fraction becomes less steep with increasing $\Psi$. The crossover takes place at $t \approx 24$ for the set of parameters considered in the present study (see, inset in Fig. 1b). There is an overall increase in the lifetime of the droplet with an increasing value of $\Psi$. This is due to the competition between the decrease in dynamic contact angle and increase in $\Psi$. Unlike, what is expected in the constant contact angle case [4], the normalised rate of mass evaporation decreases with time due to the shrinking free surface (liquid-air interface; $f(A)$ in Eq. 2) area of the droplet at the early stage. The effect of surfactant parameter is more pronounced at the late stage of evaporation, as the mole fractions of the salt and surfactant significantly increase in the droplet (Fig. 1(b)). Fig. S1 in the supplementary information shows that the corresponding decrease in the mole fraction of water nonlinearly. Both the Fig. 1 (a) and (b) shows a highly nonlinear evolution of the contact angle and the instantaneous mass-loss rates that are significantly affected by the surfactant spreading coefficient values.

Another important factor influencing the drying time of the droplets is the relative humidity. In Fig. 2(a), it can be seen that increasing RH (%) decreases the rate of decline of the normalised contact angle ($\theta/\theta_0$). This is because increasing RH decreases the evaporation rate of the pinned droplet as shown in Fig. 2(b). It is found (not shown) that the mole fraction of the water remains close to one for a longer time as the value of RH is increased and subsequently, decreases rapidly. These distinct stages are apparent in Fig. 2(b), which shows that in the early stage, the normalised instantaneous mass evaporation rate, $m_i/m_0$ is almost constant, and is followed by a rapid decrease in $m_i/m_0$ in the late stage. It can also be seen in Fig. 2(b) that the lifetime of the droplet increases significantly for RH $\geq$90%. This range of humidity corresponds to the monsoon scenario in the Indian subcontinent where the RH peaks at around 90% - 100%. We found that the droplet lifetimes are approximately equal to 32.5 s for RH $\geq$50%, 180 s for RH $=90\%$, 350 s for RH $=95\%$ and 860 s for RH $=99\%$. These values demonstrate the high sensitivity of the droplet lifetime on the relative humidity of the ambient.

To further elucidate the effect of the relative humidity, temperature, initial contact angle and $\Psi$ on the lifetime of droplets, we present the regime maps in Fig. 3(a) and (b), showing the drying time of the droplet in $T$–RH space in the two limits of initial contact angle considered in the present study, namely, $\theta_0 = 10^\circ$ and $\theta_0 = 90^\circ$, respectively for $V_0 = 10$ nL, $M = 0.154$ mol/kg and $\Psi = 20$. The corresponding region maps for $\Psi = 0$ are provided as Fig. S2 in the supplementary information. The iso-contour lines in these maps show the drying time in second. A logarithmic scale is used for the colourmap to represent the drying time of the droplet as it varies by 4 orders of magnitude. The longest drying times are observed with the combination of low $T$ and high RH (upper left corner), while the drying time progressively reduces as the humidity falls and the temperature rises for both sets of parameters considered. We also found that for a fixed initial volume of the droplet, increasing the initial contact angle increases the drying time. This follows physical intuition, as a droplet with a smaller contact angle has a greater wetting face area over which the diffusion-driven evaporation can occur. Thus, highly hydrophilic surfaces may be less susceptible to prolonged contamination in comparison to less hydrophilic surfaces, which is somewhat counter-intuitive result. Increasing the surfactant parameter, $\Psi$ also found to increase the drying time of the droplet.

Finally, in Fig. 4(a-d), we present the variations of the lifetime of the saliva droplets laden with salt and insoluble surfactants with $\Psi$ for different values of $\theta_0$, $T$, RH and $V_0$, respectively. It can be seen that increasing $\Psi$, $\theta_0$ and RH and $V_0$ and decreasing $T$ increases the lifetime of the droplet for a fixed set of other parameters. It can be seen that the drying time increases with increasing $\Psi$ for all cases. Further, the rate of increase of the drying time with $\Psi$ becomes larger as the value of the initial contact angle, relative humidity and initial droplet volume increases and the value of the ambient temperature falls. Thus the surfactant induced decrease in the rate of evaporation has a more pronounced effect in enhancing the lifetime of the droplet especially when the basal evaporation rate is slow due to the other conditions, such as high RH, low $T$, higher $\theta_0$ and higher $V_0$.

4. Concluding remarks

To summarise, in our study, a diffusion-based theoretical model is
developed which takes into account the dynamic contact angle, the salt concentration and the surfactant effect during the evaporation of sessile saliva droplets and the associated physics. The current work is an extension of Bhardwaj and Agrawal’s study [4] to consider saliva properties (such as salinity and other insoluble surfactants) and the extension of Bhardwaj and Agrawal’s study [4] to consider saliva properties (such as salinity and other insoluble surfactants) and the dynamic contact angle associated with the evaporating droplet. The effect of the surfactant was modelled by incorporating the accommodation coefficient, $\alpha$ into the diffusion-driven mass evaporation rate. The droplet is assumed to be pinned throughout its lifetime, as is expected for droplets that are saline in nature. We showed that the drying time calculated assuming saliva fluid as water and using a static contact angle significantly under-predicts the real drying time of the saliva droplets. It is observed that drying time is strongly dependent on humidity when RH $>90\%$. The regime maps in $T$–RH space show that the longest drying time is associated with the combination of low temperature and high RH values. We also found that increasing the initial contact angle and the surfactant parameter, $\Psi$ increases the drying time. The quantitative data of the lifetime of saliva droplets presented in the present study shows the importance of properly sanitising the contact prone surfaces, more so in the monsoon season and air-conditioned indoor areas.

Data availability statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Declaration of Competing Interest

We do not have any conflict of interest.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at https://doi.org/10.1016/j.icheatmasstransfer.2021.105178.

References

[1] Q. Li, X. Guan, P. Wu, X. Wang, L. Zhou, Y. Tong, R. Ren, K. Leung, E. Lau, J. Y. Wong, X. Xing, Early transmission dynamics in Wuhan, China, of novel coronavirus-infected pneumonia, N. Engl. J. Med. 382 (2020) 1199–1207.

[2] J.W. Tang, Y. Li, I. Eames, P.K.S. Chan, G.L. Ridgway, Factors involved in the aerosol transmission of infection and control of ventilation in healthcare premises, J. Hosp. Infect. 64 (2) (2006) 105–114.

[3] R. Mittal, R. Ni, J.H. Seo, The flow physics of COVID-19, J. Fluid Mech. 894 (2020) (F2–1–14).

[4] R. Bhardwaj, A. Agrawal, Likelihood of survival of coronavirus in a respiratory droplet deposited on a solid surface, Phys. Fluids 32 (6) (2020), 061704.

[5] L. Bourouiba, Turbulent gas clouds and respiratory pathogen emissions: potential implications for reducing transmission of COVID-19, JAMA 323 (18) (2020) 1837–1838.

[6] L.D. Chen, Effects of ambient temperature and humidity on droplet lifetime—a perspective of exhalation sneeze droplets with COVID-19 virus transmission, Int. J. Hyg. Environ. Health 113568 (2020).

[7] R.G. Larson, Transport and deposition patterns in drying sessile droplets, AIChE J. 60 (5) (2014) 1538–1571.

[8] E.P. Vejarano, L.C. Marr, Physico-chemical characteristics of evaporating respiratory fluid droplets, J. R. Soc. Interface 15 (139) (2018) 20170939.

[9] P. Gurrala, P. Kate, S. Balusamy, S. Banerjee, K.C. Sahu, Evaporation of ethanol-water sessile droplet of different compositions at an elevated substrate temperature, Int. J. Heat Mass Transf. 145 (2019) 118770.

[10] Y. Fukutani, D. Oregen, Y. Kitza, Y. Takata, J. Kim, K. Sefiane, Effect of ambient temperature and relative humidity on interfacial temperature during early stages of drop evaporation, Phys. Rev. E 93 (4) (2016), 043103.

[11] K. Sefiane, L. Tadrits, M. Douglas, Experimental study of evaporating water-ethanol mixture sessile drop: influence of concentration, Int. J. Heat Mass Transf. 46 (23) (2003) 4527–4534.

[12] G. Karapetian, K.C. Sahu, D.K. Matai, Evaporation of sessile droplets laden with particles and insoluble surfactants, Langmuir 32 (27) (2016) 6871–6881.

[13] V. Soslić, S. Karpińska, F. Lequien, P. Preté, T. Zemb, H. Moehwald, H. Riegler, The evaporation behavior of sessile droplets from aqueous saline solutions, Phys. Chem. Chem. Phys. 17 (34) (2015) 2296–2303.

[14] D. Brutin, V. Starov, Recent advances in droplet wetting and evaporation, Chem. Soc. Rev. 47 (2) (2018) 558–585.

[15] H. Hui, R.G. Larson, Evaporation of a sessile droplet on a substrate, J. Phys. Chem. B 106 (6) (2002) 1334–1344.

[16] R.H. Perry, D.W. Green, J.O. Maloney, Perry’s Chemical Engineers’ Handbook, 8th ed., McGrawHill Education, 2008.

[17] D.C.D. Roux, J.J. Cooper-White, Dynamics of water spreading on a glass surface, J. Colloid Interface Sci. 277 (2) (2004) 424–436.

[18] S. Chandra, C.T. Avedisian, On the collision of a droplet with a solid surface, Phys. Fluids 32 (6) (2020), 061704.

[19] V. Soulié, S. Karpińska, F. Lequien, P. Preté, T. Zemb, H. Moehwald, H. Riegler, The evaporation behavior of sessile droplets from aqueous saline solutions, Phys. Chem. Chem. Phys. 17 (34) (2015) 2296–2303.

[20] Y.-l. Hsieh, J. Thompson, A. Miller, Water wetting and retention of cotton assemblies as affected by alkaline and bleaching treatments, Text. Res. J. 66 (7) (1996) 456–464.

[21] G.P. Mantas, R.A. Young, Wetting of wood, Wood Sci. Technol. 31 (5) (1997) 339.

[22] D.C.D. Roux, J.J. Cooper-White, Dynamics of water spreading on a glass surface, J. Colloid Interface Sci. 277 (2) (2004) 424–436.

[23] E.R.I. Miles, J.P. Reid, I. Riipinen, Comparison of approaches for measuring the mass accommodation coefficient for the condensation of water and sensitivities to uncertainties in thermophysical properties, J. Phys. Chem. A 116 (44) (2012) 10810–10825.