Implications of the measurement of the \( B_s^0 \bar{B}_s^0 \) mass difference

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We analyze the significant new model independent constraints on extensions of the standard model (SM) that follow from the recent measurements of the \( B_s^0 \bar{B}_s^0 \) mass difference. The time-dependent CP asymmetry in \( B_s \to \psi \phi \), \( S_{\psi \phi} \), will be measured with good precision in the first year of LHC data taking, which will further constrain the parameter space of many extensions of the SM, in particular, next-to-minimal flavor violation. The CP asymmetry in semileptonic \( B_s \) decay, \( A_{SL}^s \), is also important to constrain these frameworks, and could give further clues to our understanding of the flavor sector in the LHC era. We point out a strong correlation between \( S_{\psi \phi} \) and \( A_{SL}^s \) in a very broad class of new physics models.

Recently the DØ \(^{[1]}\) and CDF \(^{[2]}\) collaborations reported measurements of the \( B_s^0 \bar{B}_s^0 \) mass difference

\[
17 \text{ps}^{-1} < \Delta m_s < 21 \text{ps}^{-1} \quad (90\% \text{ CL, DØ}),
\]

\[
\Delta m_s = (17.31^{+0.33}_{-0.18} \pm 0.07) \text{ps}^{-1} \quad (\text{CDF}). \tag{1}
\]

The probability of the signal being a background fluctuation is 0.2\% (5\%) for DØ (CDF). More important than the (moderate) improvements in the standard model (SM) global fit for the Cabibbo-Kobayashi-Maskawa (CKM) elements is that these measurements provide the first direct constraint on new physics (NP) contributions to the \( B_s \bar{B}_s \) mixing amplitude.

We focus below on a large class of NP models with the following features \(^{[3]}\): (I) The \( 3 \times 3 \) CKM matrix is unitary; (II) Tree-level decays are dominated by the SM contributions. These assumptions are rather mild and allow for large deviations from the SM predictions. It is therefore important to examine how present and near future experimental data constrain the parameter space of such models.

We expect NP contributions to modify the predictions for observables that are related to flavor-changing neutral current (FCNC) processes. A priori, we have no knowledge of the expected size of these contributions. However, due to the hierarchy problem, new degrees of freedom should be present near the electroweak symmetry breaking (EWSB) scale. Allowing for 10\% fine tuning in the SM Higgs potential, the new degrees of freedom which regularize the Higgs quadratic divergence from the top-loop should have masses \( m_X \sim 3 \text{ TeV} \) (see, e.g., \(^{[4]}\)). In a most generic natural theory such a particle can have tree-level non-universal couplings to the SM quarks. Thus, after integrating out \( X \), four-fermion operators of the form \((d^i \bar{d}^j)^2/m_X^2\) \((i,j = 1..3)\) are expected to be generated with order one complex coefficients. This would contribute to many well-measured processes in the \( B_0^q \) \((q = d,s)\) and \( K^0 \) systems. For instance, in \( K^0 \bar{K}^0 \) and \( B_s^0 \bar{B}_s^0 \) mixing we can parameterize the ratio between the NP and the SM short distance contributions by \( h_{K,q} e^{2\pi i \Delta m X} \). Assuming arbitrary CP violating phases, we expect the following orders of magnitudes for these parameters in the general case

\[
h_{K,q} \sim \left( \frac{4\pi v}{m^4 X \lambda^{5.3.2}} \right)^2 \sim O(10^5, 10^3), \tag{2}
\]

where \( v \) is the EWSB scale. Clearly, such huge values are excluded by many other observables, but this way of presenting the NP expectation will be useful in the following discussion. The smaller the ratio between the experimental bounds on \( h_{K,d,s} \) and \( h_{K,d,s}^{gen} \), the more disfavored this framework is.

The bounds on the above parameters prior to the \( \Delta m_s \) measurement were given in \(^{[5]}\), \( h_{K,d} \lesssim 0.6, 0.4 \), which are \( O(10^{-6}, 10^{-4}) \) times smaller than Eq. (2), while no significant bound was found on \( h_s \). The smallness of these ratios demonstrates that generic models which address the SM fine tuning problem are in great tension with indirect bounds from FCNC processes. These require that the scale of \( m_X \) should be orders of magnitude above the TeV scale.

The SM quark flavor sector is far from being generic as well. Most of the SM flavor parameters are small and hierarchical, and the flavor sector possesses an approximate \( U(3)_d \times U(2)_u \times U(2)_Q \) flavor symmetry (here \( d, u, Q \) correspond to the down and up type singlet and doublet quarks, respectively). Roughly speaking, only the top Yukawa coupling violates these approximate symmetries. Thus it is not inconceivable that the NP at \( m_X \) will share the same flavor symmetries. In this case its contributions to FCNC processes will be suppressed and Eq. (2) overestimates their size. Below we therefore assume that this is the case, and the new non-flavor-universal higher dimensional operators are invariant under these symmetries.

The special case in which these new operators are fully aligned with the SM Yukawa matrices corresponds to the minimal flavor violation (MFV) framework. Then the only sources of flavor and CP violation are due to the SM \(^{[6]}\). A more general case is when the new operators are only quasi-aligned with the SM Yukawa matrices, that is, in the basis where the new operators are flavor diagonal, the diagonalizing matrices of the Yukawa
couplings are at least as hierarchical as the CKM matrix. This constitutes next-to-minimal minimal flavor violation (NMFV)\(^3\). In this case there are new flavor and CP violating parameters, so NMFV is almost as generic as the class of models defined above by conditions (I) and (II). However, our assumption of quasi-alignment provides a useful way for “power counting” and to estimate the size of the expected NP contributions. Moreover it is also realized by many supersymmetric and non-supersymmetric models (see\(^3\) for more details), providing a powerful framework for model independent analysis.

What is the expected size of the NP contributions? Four-fermion operators are generated when the NP is integrated out at a scale of order \( \Lambda_{\text{NMFV}} \sim m_X \sim 3 \text{TeV} \). Consider, for example, the operator \((\bar{Q}_3 Q_3/\Lambda_{\text{NMFV}})^2\) defined in the interaction basis (gauge, Lorentz indices and \(O(1)\) coefficients are omitted). In the mass basis, this operator contributes to \(\Delta F = 2\) processes as \([D_L^*]_{3i} (D_L)_{ij} \bar{Q}_i Q_j / \Lambda_{\text{NMFV}}^2 \sim [V_{CKM}^*]_{ij} (V_{CKM})_{ij} \bar{Q}_i Q_j / \Lambda_{\text{NMFV}}^2\), where \(D_L\) is the rotation matrix of the down type doublet quarks. Comparing the NP contributions to the SM ones we find that within the NMFV we expect

\[ h_{K,d,s}^{\text{NMFV}} \sim O(1). \tag{3} \]

The magnitudes of \( h_{K,d,s} \) are inversely proportional to the cutoff of the theory and provide a measure of the tuning in the model. Moreover, a connection between \( \Lambda_{\text{NMFV}} \) and \( m_X \) relates this fine tuning to the one in the Higgs sector. Consequently, just as in the case of electroweak precision tests, any model of this class will be disfavored if the constraints on the \( h_{K,d,s} \) drop below the 0.1 level.

Below we focus on NP in \( \Delta F = 2\) processes, which are in general theoretically cleaner and have simpler operator structures. To constrain deviations from the SM in these processes, the tree-level observables \([V_{ub}/V_{cb}]\) and \(\gamma\) extracted from the CP asymmetry in \(B^\pm \to DK^\pm\) modes are crucial, because they are unaffected by NP. We consider in addition the following observables: the \(B^0_d\overline{B}^0_d\) mass differences, \(\Delta m_q\); CP violation in \(B^0_q\) mixing, \(A^q_{\text{SL}}\)\(^4\): the time dependent CP asymmetries in \(B^0_q\) decays, \(S_{qK}\) and \(S_{q\rho,\pi,\sigma}\); and the time dependent CP asymmetry in \(B^0_s\) decay, \(S_{\phi}\)\(^1\); the lifetime difference between the CP-even and CP-odd \(B_s\) states, \(\Delta \Gamma_{s}^{\text{CP}}\)\(^2\). (Of these, \(A^q_{\text{SL}}\) and \(S_{\phi}\) have not been measured, however, they will be important in the discussion below.)

The NP contributions to \(B^0_d\) and \(B^0_s\) mixing can be expressed in terms of four parameters, \(h_q\) and \(\sigma_q\) defined by

\[ M^q_{\text{SM}} = (1 + h_q e^{2i\sigma_q}) M^q_{\text{SM}}, \]

where \(M^q_{\text{SM}}\) is the dispersive part of the \(B^0_q\overline{B}^0_q\) mixing amplitude in the SM. (For a similar parameterization of NP in the \(K^0\) system, see\(^3\).) Then the predictions for the above observables are modified compared to the SM as follows:

\[
\begin{align*}
\Delta m_q & = \Delta m_q^{\text{SM}} |1 + h_q e^{2i\sigma_q}|, \\
S_{qK} & = \sin \left[ 2\beta + \arg \left( 1 + h_q e^{2i\sigma_q} \right) \right], \\
S_{q\phi} & = \sin \left[ 2\beta_3 - \arg \left( 1 + h_q e^{2i\sigma_q} \right) \right], \\
A^q_{\text{SL}} & = \text{Im} \left\{ \Gamma_1^{q}/\left( M^q_{\text{SM}} + 1 + h_q e^{2i\sigma_q} \right) \right\}, \\
\Delta \Gamma_{s}^{\text{CP}} & = \Delta \Gamma_{s}^{\text{SM}} \cos^2 \left[ \arg \left( 1 + h_q e^{2i\sigma_q} \right) \right].
\end{align*}
\]

Here \( \lambda \approx 0.23 \) is the Wolfenstein parameter, \( \beta_3 = \arg \left( - (V_{ts} V_{tb}^*)/(V_{us} V_{ub}^*) \right) \approx 1^\circ \) is the angle of a squashed unitarity triangle, and \(\Gamma_1^q\) is the absorptive part of the \(B^0_q\overline{B}^0_q\) mixing amplitude, which is probably not significantly affected by NP. (We neglect \(O(M^2_W/\Lambda^2_{\text{NMFV}})\) corrections due to NP contributions to SM tree-level \(\Delta F = 1\) processes; for a different approach, see\(^3\).)

Looking at Eq. (1) one notices a fundamental difference between the \(B_d\) and \(B_s\) systems. The SM contributions affecting the \(B_d\) system are related to the non-degenerate unitarity triangle. Thus the determination of \(h_d, \sigma_d\) is strongly correlated with that of the Wolfenstein parameters, \(\rho, \tilde{\eta}\). On the other hand the unitarity triangle relevant for the \(B_s\) system is nearly degenerate and therefore the determination of \(h_s, \sigma_s\) is almost independent of \(\rho, \tilde{\eta}\).

Figure\(^3\) shows the allowed \(h_s, \sigma_s\) parameter space without (left) and with (right) the measurement of \(\Delta m_s\) in Eq. (4) and the bound on \(\Delta \Gamma_{s}^{\text{CP}}\), using the CKMfitter package\(^1\). We used the constraint on the ratio

\[
\frac{\Delta m_d}{\Delta m_s} = \frac{1 + h_d e^{2i\sigma_d}}{1 + h_s e^{2i\sigma_s}} \left| \frac{V_{ts}}{V_{ts}} \right|^2 \frac{m_{B_d}}{m_{B_s}} \xi^2, \tag{5}
\]

which is theoretically cleaner than either \(\Delta m_d\) or \(\Delta m_s\). Since \(\Delta m_d\) depends on \(h_d, \sigma_d, \rho, \tilde{\eta}\), in order to produce the above plots these parameters were scanned over. We can easily see that the new measurement excludes a large part of the previously allowed parameter space. The excluded region around \(h_s = 1\) and \(\sigma_s = 90^\circ\) would give cancelling contributions to \(\Delta m_s\). The decrease in CL around \(h_s = 1\) is due to the \(\Delta \Gamma_{s}^{\text{CP}}\) constraint, which is useful at present, largely because its central value favors any deviation from the SM. After a year of LHC data, the bound from this quantity will probably be less important, because of theoretical uncertainties.

The magnitudes of the \(h_q\)'s provide a measure of how much fine tuning is required to satisfy the experimental constraints. Generically we do not expect the NP contributions to be MFV-like, i.e., aligned with the SM.

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\(^{1}\) By \(S_{\phi}\) we mean the CP asymmetry divided by \((1 - 2f_{\phi}^{\text{odd}})\) to correct for the CP-odd \(\phi\) fraction, which also equals \(-S_{\phi(t)}\).

\(^{2}\) Unless otherwise stated, the input parameters are as in\(^1\).
Thus we are interested in finding the allowed ranges of $h_i$, for $\sigma_i$ not near 0 mod $\pi/2$. The present constraints are roughly

$$h_d \lesssim 0.3, \quad h_s \lesssim 2, \quad h_K \lesssim 0.6.$$  

(6)

Let us now discuss some implications of the above results. Equation (6) shows that at present none of the bounds on the NP parameters have reached the 0.1 level, so NMFV survives the current tests. It is then interesting to ask which future measurements will be most important to verify or disfavor the NMFV framework. The constraints on $h_{d,K}$, even though they underwent significant improvements in the last few years due to new SM tree-level measurements [11], are now limited by the statistical errors in the measurements of $\gamma$ (and effectively $\alpha$) and the hadronic parameters in the determination of $|V_{ub}|$ from semileptonic decays and $|V_{td}|$ from $\Delta m_d$. The improvements in these constraints will be incremental, as they depend on the integrated luminosities at the $B$ factories and on progress in lattice QCD. The constraint from $\varepsilon_K$ on the $K$ system is also dominated by hadronic uncertainties. At present, the bound on $h_s$ is weaker than that on $h_d$, since only one measurement, $\Delta m_s$, constrains it, and the hadronic uncertainties are comparable.

However the $B_s$ system is exceptional because a measurement of $S_{\psi\phi}$ (or a strong bound on it) would provide a very sensitive test of NMFV, which is neither obscured by hadronic uncertainties nor by uncertainties in the CKM parameters. In the SM $S_{\psi\phi}$ is suppressed by $\lambda^2$ (the SM CKM fit gives $\sin2\beta S = 0.038 \pm 0.003$), whereas Eq. (4) implies

$$S_{\psi\phi} = -\frac{h_s \sin(2\sigma_s)}{1 + h_s e^{2i\sigma_s}} + \frac{\sin(2\beta_s) + \sin(2\beta_s)}{1 + h_s e^{2i\sigma_s}},$$  

(7)

where we set $\cos2\beta_s$ to unity. Thus when the sensitivity of the measurement of $S_{\psi\phi}$ reaches the SM level, it will provide us with a strong test of NMFV. The precision that will be achieved in forthcoming experiments depends on the value of $\Delta m_s$, but since we now know $\Delta m_s$, we can use the LHC projections for the SM case. LHCb expects to reach $\sigma(S_{\psi\phi}) \approx 0.03$ with the first year (2 fb$^{-1}$) data [12] (in several years the uncertainty may be reduced to 0.01). Figure 2 shows the resulting constraint on $h_s, \sigma_s$, assuming an experimental measurement $S_{\psi\phi} = 0.04 \pm 0.03$. This plot demonstrates that already with one year of data the bound on $h_s$ will be better than 0.1, which will severely constrain the NMFV class of models. Even initial data on $S_{\psi\phi}$ at the Tevatron may constrain new physics in $B_s$ mixing comparable to similar bounds on $h_d, \sigma_d$ in the $B_d$ sector.

Another sensitive probe of this class of models is the $CP$ asymmetry in semileptonic $B_s$ decays, $A_{SL}$. In the SM it is unobservably small, because the short distance contributions are much larger than the long distance part, $|\Gamma_{12}/M_{12}^2| \propto m_s^2/m_t^2$, and the two contributions are highly aligned, $\arg(\Gamma_{12}/M_{12}^2) \propto (m_s^2/m_t^2) \sin2\beta_s$ [7]. Given the new $\Delta m_s$ result, we know that even in the presence of NP the first suppression factor can only be moderately affected, while the second one can be signif-
icantly enhanced in the presence of new CP violating phases. Figure 3 shows the allowed range of $A_{\text{SL}}^s$, taking into account the new constraint from $\Delta m_s$. We find

$$A_{\text{SL}}^s < 0.01,$$

which extends three order of magnitude above the SM prediction [13]. In particular, $|A_{\text{SL}}^s| > |A_{\text{SL}}^0|$ is possible, contrary to the SM, in which $|A_{\text{SL}}^s/A_{\text{SL}}^0| \approx |V_{td}/V_{ts}|^2$. This demonstrates that while the constraint from the $\Delta m_s$ measurement is of great importance, it still leaves plenty of room for deviations from the SM within NMFV.

Finally we point out that $A_{\text{SL}}^s$ and $S_{\psi\phi}$ are highly correlated in the region in which $h_s, \sigma_s \gg \beta_s$ and $h_s$ is moderate. Defining $2\theta_s = \arg(1 + h_s e^{2i\sigma_s})$, we have $S_{\psi\phi} = \sin(2\beta_s - 2\theta_s)$, so $A_{\text{SL}}^s$ can be written as

$$A_{\text{SL}}^s = \left[ \frac{\Gamma_{12}^{\text{SM}}}{M_{12}^{\text{SM}}} \right] \sin(2\theta_s) + \mathcal{O}(h_s^2, m_{\tau}^2/m_S^2).$$

Figure 4 shows $A_{\text{SL}}^s$ as a function of $S_{\psi\phi}$, taking into account the constraint from $\Delta m_s$ [without neglecting the $\mathcal{O}(h_s^2, m_{\tau}^2/m_S^2)$ terms]. As explained above, the two observables are strongly correlated. Deviation from this prediction would provide a clear indication of new physics beyond the generic framework defined by (1) and (II).

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