ROTATION OF THE OPTICAL POLARIZATION ANGLE ASSOCIATED WITH THE 2008 $\gamma$-RAY FLARE OF BLAZAR W COMAE

MARCO SORIA$^1$, ERIKA BENÍTEZ$^1$, DAVID HIRIART$^2$, JOSÉ M. LÓPEZ$^2$, JOSÉ I. CABRERA$^1$, AND RAÚL MÚJICA$^3$

$^1$ Instituto de Astronomía, Universidad Nacional Autónoma de México, Apdo. 70-264, Mexico D.F., 04510, Mexico
$^2$ Instituto de Astronomía, Universidad Nacional Autónoma de México, Apdo. 810, Ensenada, B.C., 22800, Mexico
$^3$ Instituto Nacional de Astro física, Óptica y Electrónica, Apdo. Postal 51 y 216, 72000 Tonantzintla, Puebla, Mexico

Received 2013 September 17; accepted 2014 July 31; published 2014 September 23

ABSTRACT

An $R$-band photopolarimetric variability analysis of the TeV bright blazar W Comae between 2008 February 28 and 2013 May 17 is presented. The source showed a gradual tendency to decrease its mean flux level with a total change of 3 mJy. A maximum and minimum brightness states in the $R$ band of $14.25 \pm 0.04$ and $16.52 \pm 0.1$ mag, respectively, were observed, corresponding to a maximum variation of $\Delta F = 5.40$ mJy. We estimated a minimum variability timescale of $\Delta t = 3.3$ days. A maximum polarization degree $P = 33.8% \pm 1.6\%$, with a maximum variation of $\Delta P = 33.2\%$, was found. One of our main results is the detection of a large rotation of the polarization angle from $78^\circ$ to $315^\circ (\Delta \theta \sim 237^\circ)$ that coincides in time with the $\gamma$-ray flare observed in 2008 June. This result indicates that both optical and $\gamma$-ray emission regions could be co-spatial. During this flare, a correlation between the $R$-band flux and polarization degree was found with a correlation coefficient of $r_{F-P} = 0.93 \pm 0.11$. From the Stokes parameters, we infer the existence of two optically thin synchrotron components that contribute to the polarized flux. One of them is stable with a constant polarization degree of $11\%$. Assuming a shock-in-jet model during the 2008 flare, we estimated a maximum Doppler factor $\delta_D \sim 27$ and a minimum of $\delta_D \sim 16$; a minimum viewing angle of the jet $\sim 2^\circ$; and a magnetic field $B \sim 0.12$ G.

Key words: BL Lacertae objects: individual (ON231, W Comae) – galaxies: jets – galaxies: photometry – polarization

Online-only material: machine-readable table

1. INTRODUCTION

The study of the blazar phenomenon has been one of the major topics of study on the active galactic nuclei (AGNs) family because of their extreme properties. They show strong flux variability, superluminal motion, and a non-thermal continuum extending from radio to TeV $\gamma$-ray regions (e.g., Jorstad et al. 2005; Abdo et al. 2010a; Marscher et al. 2010; Agudo et al. 2013). Blazars are radio-loud AGNs and consist of BL Lacertae objects and flat spectrum radio sources (Angel & Stockman 2013). Blazars are radio-loud AGNs and consist of BL Lacertae objects and flat spectrum radio sources (Angel & Stockman 2013). These properties are explained through the idea that blazars are objects with a very small viewing angle, i.e., the emission produced by the relativistic jet is aligned very close to the observer’s line of sight (e.g., Blandford & Rees 1978; Hovatta et al. 2009). In recent years, it has been well established that the non-thermal continuum emission in blazars shows two broad low-frequency and high-frequency components in their spectral energy distribution (SED). In the case of the BL Lac objects, this empirical property conforms the base for classifying them accordingly to the location of the first peak, known as the synchrotron peak, in the SED (Padovani & Giovannini 1995; Nieppola et al. 2006; Costamante 2012). Commonly, low-frequency peak BL Lac objects (LBL) have their synchrotron peak, $v_{\text{peak}} \sim 10^{14} \text{Hz}$; intermediate-frequency peak BL Lac objects (IBL) in the range $10^{14} < v_{\text{peak}} < 10^{15} \text{Hz}$; and high-frequency peak BL Lac objects (HBL) have $v_{\text{peak}} > 10^{15} \text{Hz}$.

The blazar W Comae at $z = 0.102$ (also known as 1219+285 or ON 231) was discovered as a radio source by Browne (1971). Very long baseline interferometry (VLBI) observations of W Comae revealed a complex jet that extends toward the east at $\theta \sim 100^\circ$ (Gabuzda et al. 1992, 1994). Also, it was found that the jet shows superluminal components with strong polarization. The polarized emission components are found to be both aligned with and transverse to the local jet direction in different jet components (Gabuzda & Cawthorne 1996).

The optical historical light curve of W Comae shows variations at all scales, from days and weeks to months and years (see, e.g., Liu et al. 1995; Belokon et al. 2000; Tosti et al. 1998). Also, it has shown rapid variations on scales of hours (Babadzhanov & Belokon’ 2002). Tosti et al. (1998) observed the highest brightness value ever observed for W Comae since 1940, reaching a maximum of $B = 14.2$ mag in 1997 January. Later, Massaro et al. (1999) reported a very strong flare of W Comae when the object reached a historical maximum of $R \sim 12.2$ mag in 1998 April 23. Optical polarization of W Comae was also reported in Massaro et al. (1999). Their multi-band optical observations were done just before and during its brightest phase (1998 April 17–25). During the brightest state, the polarization was low in the $UBV$ filters ($\sim 2\%$ to $4\%$) with less than $0.4\%$ in the $R_c$ and $I_c$ filters.

The $\gamma$-ray emission of W Comae has been detected by the Energetic Gamma Ray Experiment Telescope on board of the Compton Gamma-Ray Observatory in the 100 MeV to 10 GeV band (Hartman et al. 1999). BeppoSAX data analysis of W Comae given by Tagliaferri et al. (2000) demonstrates that this source is an IBL source. This blazar was considered as a very interesting target for the very high energy (VHE) observatories due to the possibility of being a $\gamma$-ray source that could be detected by Cerenkov telescopes such as the High Energy Stereoscopic System, the MAGIC telescopes, and the Very Energetic Radiation Imaging Telescope Array System (VERITAS; see, e.g., Böttcher et al. 2002). This prediction was confirmed later, when W Comae was discovered to be a $\gamma$-ray emitter at VHE by VERITAS in 2008 March 15 (see Acciari
et al. 2008). Thus, W Comae is the first IBL detected at VHE. A subsequent multi-wavelength campaign on this object was coordinated during a major γ-ray flare in 2008 June (Acciari et al. 2009). A very high γ-ray signal was detected by VERITAS in 2008 June 8 that was brighter, by a factor of three, than the previous emission detected in 2008 March.

In this paper, we report the results of the photopolarimetric monitoring of the TeV-blazar W Comae carried out from 2008 February to 2013 May. Our main goal is to establish the long-term optical variability properties of the polarized emission in the $R$ band. The variability of the Stokes parameters obtained from our observations is analyzed in terms of a two-component model. Estimations of some of the physical parameters that are known to be associated with the kinematics of the relativistic jet are obtained. One of our main results is the detection of a large rotation of the electric vector position angle (EVPA) that coincides with the time of occurrence of the major flare observed in γ-rays in 2008 June 8.

This paper is organized as follows. Section 2 presents a description of our observations and our main observational results. Polarimetric properties are analyzed in Section 3. Results are discussed in Section 4. Finally, in Section 5, we show our conclusions. Throughout this paper, we use a standard cosmology with $H_0 = 71$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_m = 0.27$, and $\Omega_\Lambda = 0.73$.

2. OBSERVATIONS AND RESULTS

The observations were carried out with the 0.84 m f/15 Ritchey–Chretien telescope at the Observatorio Astronómico Nacional of San Pedro Mártir (OAN–SPM) in Baja California, Mexico and the instrument POLIMA. The differential $R$-band magnitudes of W Comae were calculated using the standard star $\gamma$ coincides with the time of occurrence of the major flare observed in 2008 June (Acciari et al. 2009). A very high $\gamma$-ray signal was detected by VERITAS in 2008 June 8 that was brighter, by a factor of three, than the previous emission detected in 2008 March.

In this paper, we report the results of the photopolarimetric monitoring of the TeV-blazar W Comae carried out from 2008 February to 2013 May. Our main goal is to establish the long-term optical variability properties of the polarized emission in the $R$ band. The variability of the Stokes parameters obtained from our observations is analyzed in terms of a two-component model. Estimations of some of the physical parameters that are known to be associated with the kinematics of the relativistic jet are obtained. One of our main results is the detection of a large rotation of the electric vector position angle (EVPA) that coincides with the time of occurrence of the major flare observed in $\gamma$-rays in 2008 June 8.

This paper is organized as follows. Section 2 presents a description of our observations and our main observational results. Polarimetric properties are analyzed in Section 3. Results are discussed in Section 4. Finally, in Section 5, we show our conclusions. Throughout this paper, we use a standard cosmology with $H_0 = 71$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_m = 0.27$, and $\Omega_\Lambda = 0.73$.

2.1. Global Variability Properties

W Comae was observed between 2008 February 28 and 2013 May 17. During this period, 32 observing runs of 7 nights per run were carried out around the new moon phase; in total, we collected 141 data points. The observational results are presented in Table 1, where Column 1 is the observation cycle (see explanation in next paragraph); Columns 2 and 3 give the Gregorian and Julian Dates of the observation, respectively; Columns 4 and 5 give the polarization degree and its error, respectively; Columns 6 and 7 give the polarization position angle, a $\chi^2$ test was carried out.

| Cycle | JD 2,450,000+ | $p$ (%) | $\epsilon_p$ (%) | $\theta$ (°) | $\epsilon_\theta$ (°) | $R$ (mag) | $\epsilon_R$ (mag) | Flux (mJy) | $\epsilon_{\text{Flux}}$ (mJy) |
|-------|--------------|--------|-----------------|-------------|-----------------|--------|----------------|----------|------------------|
| I     | 2008 Feb 28  | 4524.9351 | 11.0 | 0.3 | 96 | 01 | 14.80 | 0.04 | 3.71 | 0.07 |
|       | 2008 Feb 29  | 4525.9038 | 10.4 | 0.2 | 93 | 01 | 14.53 | 0.04 | 4.75 | 0.08 |
|       | 2008 Mar 1   | 4526.9429 | 13.7 | 0.2 | 82 | 01 | 14.45 | 0.04 | 5.11 | 0.09 |
|       | 2008 Mar 9   | 4534.8984 | 11.1 | 0.4 | 83 | 01 | 14.59 | 0.04 | 4.50 | 0.08 |
|       | 2008 Mar 10  | 4535.8706 | 10.8 | 0.3 | 78 | 01 | 14.57 | 0.04 | 4.60 | 0.08 |
|       | 2008 Mar 11  | 4536.8555 | 9.6 | 0.4 | 88 | 01 | 14.70 | 0.04 | 4.04 | 0.07 |
|       | 2008 Mar 12  | 4537.9004 | 6.3 | 0.3 | 103 | 02 | 14.47 | 0.04 | 5.00 | 0.08 |
|       | 2008 Mar 15  | 4540.8467 | 12.8 | 0.3 | 97 | 01 | 14.53 | 0.04 | 4.75 | 0.08 |
|       | 2008 May 5   | 4591.8618 | 5.9 | 0.9 | 113 | 03 | 14.42 | 0.04 | 5.27 | 0.09 |
|       | 2008 May 7   | 4593.7847 | 8.2 | 1.5 | 105 | 04 | 14.53 | 0.05 | 4.75 | 0.09 |

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)
Figure 1. Photopolarimetric light curves of W Comae from 2008 February 28 to 2013 May 17: (from top to bottom) $R$-band flux, $F$ (mJy); magnitude, $R$ (mag); percentage of linear polarization in the $R$ band, $p$ ($\%$); and orientation of the EVPA, $\theta$ (°). Vertical dashed lines separate the monitoring period into Cycles I to VI. Associated errors are presented in Table 1. $R$-band magnitudes and fluxes have been corrected for the host galaxy contribution (see text for explanation). The $R$-band light curve (top panel) shows a slow decreasing of the mean level flux (dashed line) with a fall of $\sim$3 mJy in 5.2 yr. The amplitude of the variations $Y (\%)$ was estimated using flux densities instead of magnitude differences following Heidt & Wagner (1996),

$$Y (\%) = \frac{100}{\langle S \rangle} \sqrt{\langle S_{\text{max}} - S_{\text{min}} \rangle^2 - 2\sigma_c^2},$$

where $S_{\text{max}}$ and $S_{\text{min}}$ are the maximum and minimum values of the flux density, respectively. $\langle S \rangle$ is the mean value and $\sigma_c^2 = \sigma_{\text{max}}^2 + \sigma_{\text{min}}^2$. The variability is described by the fluctuation index $\mu$ defined by

$$\mu = 100 \frac{\sigma S}{\langle S \rangle} \%,$$

and the fractional variability index of the source $\mathcal{F}$ obtained from the individual nights:

$$\mathcal{F} = \frac{S_{\text{max}} - S_{\text{min}}}{S_{\text{max}} + S_{\text{min}}}.$$

We have estimated the minimum flux variability timescale using the definition proposed by Burbidge et al. (1974):

$$\tau = \frac{dt}{\ln(F_1/F_2)} \ ,$$

where $\tau$ is the time interval between flux measurements $F_1$ and $F_2$, with $F_1 > F_2$. We have calculated all possible timescales $\tau_{ij}$ for any pair of observations for which $|F_i - F_j| > \sigma_{F_i} + \sigma_{F_j}$. The minimum timescale is obtained when

$$\tau_{\text{var}} = \min\{\tau_{ij}, \nu\},$$

where $i = 1, ..., N - 1; j = i + 1, ..., N$, and $N$ is the number of observations. The uncertainties associated to $\tau_{\text{var}}$ were obtained through the errors in the flux measurements.

Table 2 shows the results obtained from the statistical analysis: Column 1 gives the corresponding cycle; Column 2 the variable parameters; Columns from 3 to 10 present, for each of the four parameters, the average, the maximum and minimum observed value, the maximum variation $\Delta_{\text{max}}$, the variability amplitude $Y (\%)$, the variability index $\mu (\%)$, the variability fraction $\mathcal{F}$, and the statistic $\chi^2$, respectively. We have estimated the minimum flux variability timescale of $\tau_{\text{var}} = 3.3 \pm 0.3$ d.

2.2. Photometric Variability

Considering the entire data set, a brightness maximum of $R = 14.25$ mag was observed in 2008 June 4 and a brightness minimum of $R = 16.52$ mag in 2013 May 17. A variation of $\Delta m_R = 2.27$ mag (5.40 mJy) in $\Delta t = 1905$ d ($\sim$5.2 yr) is found (see Table 2). During our monitoring period, the source showed a maximum brightness variation in timescales from months to years. There can be noticed a tendency of a slow decreasing brightness after each flare episode, which is shown in Figure 1. In this figure, a fall of $\sim$3 mJy in $\sim$5.2 yr, superimposed on rapid brightness variations with timescales of months and days, can be seen. The time between peak brightness maxima is $\sim$0.9–1.0 yr.

The most important photometric results are found in Cycles I, V, and VI (see Table 2). In Cycle I, W Comae shows a maximum flux of $6.16 \pm 0.10$ mJy in 2008 June 4. This flare lasted $\sim$ two months. A minimum flux of $3.71 \pm 0.07$ mJy is observed in 2008 February 28. The flux changed 2.45 mJy in 97 days. We want to point out here that all photometric $R$-band data
In this long-term flare, there are two superimposed short-term cycles corresponding to a long-term flare (2008 June 3–7 flare). In Cycle III, its value is \( \Delta P = 0.09 \) during the rise of the flare. In Cycle VI, its value is \( \Delta P = 0.93 \) during the fall of the flare. This result points out that both the flux and the polarization degree show a tendency to be anti correlated in periods of time \( \sim \) weeks to months. On the other hand, a positive correlation of \( r_{F-P} = 0.93 \) was found during Cycle I (2008 June 3–7 flare). In general, the polarization degree showed a random variability behavior, with a maximum and a minimum of \( (33.8 \pm 1.6)\% \) (2013 May 12) and \( (0.6 \pm 1.0)\% \) (2008 July 9), respectively. The maximum variability observed was \( \Delta P = 33.2\% \), in \( \Delta t = 1768 \) days \( \sim 4.8 \) yr. It is interesting to notice that the maximum value of the polarization degree occurred in Cycle VI, when the brightness was at its minimum.

The maximum and minimum polarization degrees for each cycle are shown in the Table 2. In Cycle I, the maximum variability observed of the polarization degree is \( \Delta P = 13.1\% \) in \( \Delta t = 130 \) days; in Cycle II, \( \Delta P = 14.7\% \) in \( \Delta t = 6 \) days; in Cycle III, \( \Delta P = 7.8\% \) in \( \Delta t = 57 \) days; in Cycle IV, \( \Delta P = 14.5\% \) in \( \Delta t = 92 \) days; in Cycle V, \( \Delta P = 17.8\% \) in \( \Delta t = 39 \) days; and in Cycle VI, \( \Delta P = 28.5\% \) in \( \Delta t = 55 \) days.

2.3.2. Position Angle Variability

In general, our data do not show a clear correlation between the polarization angle and the R-band flux (see bottom panel of Figure 2). Rather, after the large rotation observed during

| Cycle (1) | Parameter (2) | Average (3) | Max (4) | Min (5) | \( \Delta_{\text{max}} \) (6) | \( Y(\%) \) (7) | \( \mu(\%) \) (8) | \( J \) (9) | \( \chi^2 \) (10) |
|----------|--------------|-------------|---------|--------|----------------|----------------|---------------|---------|-------------|
| All      | \( R(\text{mag}) \) | 15.06 ± 0.048 | 16.52 | 14.25 | 2.27 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 3.17 ± 0.126 | 6.16 | 0.76 | 5.40 | 170.1 | 39.6 | 0.78 | 62705.7 |
|          | \( P(\%) \) | 12.52 ± 0.063 | 33.82 | 0.59 | 33.23 | 264.6 | 50.7 | 0.97 | 4404.1 |
|          | \( \theta(\%) \) | 72.49 ± 33.40 | 314.59 | 5.54 | 309.05 | 425.1 | 46.1 | 0.97 | 13546.4 |
| I        | \( R(\text{mag}) \) | 14.50 ± 0.015 | 14.80 | 14.25 | 0.55 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 4.91 ± 0.066 | 6.16 | 3.71 | 2.45 | 49.8 | 13.5 | 0.25 | 1119.3 |
|          | \( P(\%) \) | 7.86 ± 0.038 | 13.70 | 0.59 | 13.11 | 165.7 | 48.9 | 0.92 | 979.7 |
|          | \( \theta(\%) \) | 124.86 ± 58.80 | 314.59 | 77.99 | 236.60 | 188.6 | 47.1 | 0.60 | 1858.6 |
| II       | \( R(\text{mag}) \) | 14.50 ± 0.007 | 14.67 | 14.36 | 0.31 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 4.90 ± 0.031 | 5.55 | 4.18 | 1.38 | 27.9 | 6.4 | 0.14 | 212.3 |
|          | \( P(\%) \) | 9.29 ± 0.032 | 19.91 | 5.20 | 14.71 | 157.6 | 35.7 | 0.59 | 291.2 |
|          | \( \theta(\%) \) | 74.90 ± 13.33 | 116.57 | 55.02 | 61.55 | 80.2 | 17.8 | 0.36 | 211.4 |
| III      | \( R(\text{mag}) \) | 14.97 ± 0.018 | 15.30 | 14.72 | 0.58 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 3.21 ± 0.052 | 4.00 | 2.35 | 1.65 | 51.4 | 16.3 | 0.26 | 1262.7 |
|          | \( P(\%) \) | 15.75 ± 0.027 | 18.71 | 10.93 | 7.78 | 48.5 | 15.0 | 0.26 | 205.0 |
|          | \( \theta(\%) \) | 61.16 ± 07.98 | 77.96 | 49.76 | 28.20 | 45.8 | 13.1 | 0.22 | 431.3 |
| IV       | \( R(\text{mag}) \) | 15.32 ± 0.020 | 15.56 | 14.76 | 0.80 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 2.34 ± 0.048 | 3.83 | 1.83 | 2.00 | 85.2 | 20.5 | 0.35 | 1885.4 |
|          | \( P(\%) \) | 8.65 ± 0.041 | 16.93 | 2.41 | 14.52 | 165.3 | 47.5 | 0.75 | 377.0 |
|          | \( \theta(\%) \) | 53.21 ± 22.47 | 113.54 | 5.54 | 108.00 | 200.9 | 42.2 | 0.91 | 1970.1 |
| V        | \( R(\text{mag}) \) | 15.08 ± 0.029 | 15.68 | 14.53 | 1.15 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 2.96 ± 0.073 | 4.75 | 1.65 | 3.09 | 104.6 | 24.6 | 0.48 | 6645.1 |
|          | \( P(\%) \) | 14.31 ± 0.045 | 20.75 | 2.89 | 17.86 | 123.3 | 28.3 | 0.76 | 312.8 |
|          | \( \theta(\%) \) | 67.91 ± 10.66 | 88.82 | 42.14 | 46.68 | 68.0 | 15.7 | 0.36 | 472.1 |
| VI       | \( R(\text{mag}) \) | 15.68 ± 0.051 | 16.52 | 14.84 | 1.68 | ... | ... | ... | ... |
|          | \( F(\text{mJy}) \) | 1.82 ± 0.085 | 3.55 | 0.76 | 2.80 | 153.7 | 46.7 | 0.65 | 7055.9 |
|          | \( P(\%) \) | 17.86 ± 10.05 | 33.82 | 5.32 | 28.50 | 158.9 | 56.3 | 0.73 | 1136.3 |
|          | \( \theta(\%) \) | 66.22 ± 17.77 | 121.93 | 44.43 | 77.50 | 116.9 | 26.8 | 0.47 | 2393.4 |

Note. There are no statistics \( Y, \mu, J \), and \( \chi^2 \) for the magnitude due to its logarithmic character.
Figure 2. Correlations between photopolarimetric observations of W Comae for all data. Top: correlation between the $R$-band flux and the polarization degree. Bottom: correlation between the $R$-band flux and EVPA. The dotted line at the bottom panel shows the preferred EVPA of 65°.

Table 3

| Cycle | Relation Parameters | $r$  | $r$-student | Confidence Level |
|-------|---------------------|-----|-------------|-----------------|
|       |                     | (3) | (4)         | (5)             |
| All   | $F - p$             | -0.44 $\pm$ 0.02 | 5.72  | no             |
|       | $F - \theta$        | 0.47 $\pm$ 0.01  | 6.25  | no             |
|       | $p - \theta$        | -0.34 $\pm$ 0.02 | 4.17  | no             |
| I     | $F - p$             | 0.93 $\pm$ 0.11  | 5.04  | yes            |
|       | $F - \theta$        | -0.90 $\pm$ 0.11 | 4.15  | yes            |
|       | $p - \theta$        | -0.92 $\pm$ 0.15 | 4.60  | yes            |
| II    | $F - p$             | -0.88 $\pm$ 0.24 | 4.60  | yes            |
|       | $F - \theta$        | -0.71 $\pm$ 0.24 | 2.47  | no             |
|       | $p - \theta$        | 0.53 $\pm$ 0.27  | 1.54  | no             |
| III   | $F - p$             | -0.82 $\pm$ 0.09 | 4.24  | yes            |
|       | $F - \theta$        | 0.81 $\pm$ 0.08  | 4.10  | yes            |
|       | $p - \theta$        | -0.68 $\pm$ 0.04 | 2.81  | no             |
| IV    | $F - p$             | -0.75 $\pm$ 0.07 | 3.18  | no             |
|       | $F - \theta$        | -0.53 $\pm$ 0.16 | 1.75  | no             |
|       | $p - \theta$        | 0.25 $\pm$ 0.14  | 0.75  | no             |
| V     | $F - p$             | 0.16 $\pm$ 0.08  | 0.98  | no             |
|       | $F - \theta$        | -0.39 $\pm$ 0.08 | 2.63  | no             |
|       | $p - \theta$        | -0.22 $\pm$ 0.13 | 1.37  | no             |
| VI    | $F - p$             | -0.89 $\pm$ 0.04 | 7.68  | yes            |
|       | $F - \theta$        | 0.41 $\pm$ 0.05  | 1.81  | no             |
|       | $p - \theta$        | -0.61 $\pm$ 0.05 | 3.07  | no             |

Notes. Pearson’s correlation coefficient, $r$, between the observed parameters: flux ($F$); polarization degree ($p$); and polarization angle ($\theta$). In order to verify the validity of the correlation to 99%, we applied a $t$-student test.

Cycle I, the polarization angle presents a preferential position of ~65° (see Section 3) with maximum variations of $\Delta \theta \sim 54^\circ$ (see bottom panel of Figure 1).

In Cycle I, a gradual rotation of the EVPA of 78° (2008 March 10) to 315° (2008 July 9) is observed. This corresponds to a total rotation of ~237° in a period of 121 days (giving an average rate of rotation of ~2° per day). Figure 3 shows this rotation in the Stokes plane. For more clarity only the more representative points are shown.

In Cycles II to VI, our data show that EVPA have the preferential value mentioned above with mean variations rate $\sim 1\pm 2$° per day. In Cycle IV, the EVPA reach a maximum value of 114° while the polarization degree is at its minimum value of 2.4%. And the other way around, when the EVPA shows its minimum value of 6°, the polarization degree shows its maximum value of 16.9%.

3. POLARIMETRIC ANALYSIS

From our observations, we have found that W Comae shows, in general, a random polarimetric behavior. This has been explained as due to the presence of one or more variable polarization components overlaid on a stable one. To identify the presence of a stable polarized component, we have used the method suggested by Jones et al. (1985). In this work, the authors proposed that if the observed average values ($\langle Q \rangle$, $\langle U \rangle$) in the absolute Stokes parameters plane $Q$–$U$ deviate significantly from the origin, then a stable polarization component is present. From our data, the derived average values of the absolute Stokes parameters are $\langle Q \rangle = -0.22 \pm 0.02$ mJy and $\langle U \rangle = 0.21 \pm 0.03$ mJy. These average values correspond to a stable component with constant polarization degree $P_e = 10.7\% \pm 0.8\%$ and
polarization angle \( \Theta_c = 65^\circ \pm 2^\circ \). The constant polarization degree has a dispersion \( \sigma_{P_c} = 6.4\% \).

We estimated the polarization variable component parameters, we looked for a possible linear relation between \( Q \) versus \( I \) and \( U \) versus \( I \) for the six relevant cycles (see, Hagen-Thorn et al. 2008). For Cycle IV, no linear correlation between these parameters was found; rather, they appear to be randomly related. In contrast, for Cycles I, II, III, V, and VI, our data show a linear tendency between these parameters. We made a least-square fit to the data in order to find the slopes and the linear correlation coefficients, \( r_{QI} \) and \( r_{UI} \). Figure 4 shows this linear correlations between Stokes parameters for Cycles I and VI. The correlation coefficients for these parameters are given in Table 4 where Columns 2 to 7 give the parameters, \( q_{\text{var}} \), \( r_{QI} \), \( u_{\text{var}} \), \( r_{UI} \), \( p_{\text{var}} \) and \( \theta_{\text{var}} \), respectively. The maximum polarization degree found for the variable component is \( p_{\text{max}}^{\text{var}} = (40.1 \pm 5.1)\% \), with a polarization angle \( \theta_{\text{var}} = 116^\circ \pm 7^\circ \), corresponding to Cycle I.

### 3.1. The Two-component Model

From the above results, we infer the presence of a stable component that we assume associated with the relativistic jet and also a variable component that can be related to the propagation of a shock. Therefore, the observed polarization would be the result of the overlap of these two optically thin synchrotron components.

Assuming that there are two polarimetric components in W Comae, we have used Equations (1) and (2) given in Holmes et al. (1984) and derive the following equations for the parameters associated to the polarized variable component:

\[
\begin{align*}
    p_{\text{var}}^2 &= p_{\text{cons}}^2 + p^2(1 + I_{v/c})^2 - 2 p p_{\text{cons}} (1 + I_{v/c}) \cos 2(\Theta_c - \theta) \frac{I_{v/c}}{I_{v/c}} ,
    \\
    \tan 2\theta_{\text{var}} &= \frac{p(1 + I_{v/c}) \sin 2\theta - p_{\text{cons}} \sin 2\Theta_c}{p(1 + I_{v/c}) \cos 2\theta - p_{\text{cons}} \cos 2\Theta_c} .
\end{align*}
\]

where \( I_{v/c} \) is the flux ratio between the variable to the constant component, and \( p \) and \( \theta \) are observed polarimetric parameters.

This system of equations has five free parameters: \( p_{\text{cons}}, \Theta_c, p, \theta, \) and \( I_{v/c} \). The system can be resolve if \( p_{\text{cons}} \) and \( \Theta_c \) correspond to \( P_{\text{cons}} \) and \( \Theta_c \) previously obtained in Section 3.

To obtain \( I_{v/c} \), we maximize Equation (7) with respect to \( \theta \). From our observations, \( p_{\text{var}} \) reaches maximum values

---

**Figure 4.** Left: linear correlation between the Stokes parameter \( Q \) vs. \( I \) (top panel), and \( U \) vs. \( I \) (bottom panel) for Cycle I. The correlation coefficients for this cycle are \( r_{QI} = 0.80 \) and \( r_{UI} = 0.96 \), and the slopes \( m_{QI} = -0.24 \pm 0.07 \) and \( m_{UI} = -0.32 \pm 0.04 \), respectively. Right: linear correlation between the Stokes parameter \( Q \) vs. \( I \) (top panel), and \( U \) vs. \( I \) (bottom panel) for Cycle VI. The correlation coefficients for this cycle are \( r_{QI} = 0.51 \) and \( r_{UI} = 0.89 \), and the slopes \( m_{QI} = 0.10 \pm 0.07 \) and \( m_{UI} = 0.33 \pm 0.07 \), respectively.
where the blazar presents a minimum activity state, we find that polarized component. For more details, see Section 3.1.

The large rotation of EVPA can be interpreted as due to an increase in the value of the EVPA from 78°, just a day before the huge gamma-ray flare. With these values we estimate $I_{\nu/c} = 0.57 \pm 0.07$. Applying the same procedure in Cycle VI, where the blazar presents a minimum activity state, we find that $I_{\nu/c} = 3.98 \pm 0.32$.

The values of $p_{\text{var}}$ and $\theta_{\text{var}}$ are shown in Figure 5, where the observed polarization $p$ is the combination of the two polarization components (stable plus variable). For Cycle I (high activity state), it can be seen that the variable polarization component $p_{\text{var}}$ shows a similar variability behavior as the observed flux in the $R$-band. We previously assumed that this variable polarization component is associated with the propagation of shocks along the jet. In the same figure, we show the results for Cycle VI (low activity state) where the observed polarization $p$ and the variable polarization component $p_{\text{var}}$ show a similar variability behavior. It is interesting to note that $\theta_{\text{var}}$ follows the observed EVPA variations in both cycles.

4. DISCUSSION

In this work, we have inferred the presence of two components to explain the optical polarization variability. The minimum variability scale of approximately three days was found and it is superimposed on a longer-term flare that lasts approximately three months (these long-term flares appeared separated by ~0.9 yr). The variability timescales found in this work are in agreement with previous studies (Tosti et al. 1998; Massaro et al. 1999).

In 2008 June 8, a strong outburst of very high energy gamma-ray emission above 200 GeV was detected with VERITAS in W Comae with a significance of 10.3 (Acciari et al. 2009). Data from our monitoring for 2008 June 4–7 show an increase in the $R$-band flux. Unfortunately, due to bad weather we could not obtain data for June 8, when the maximum brightness was observed in the $\gamma$-rays. However, our data show a gradual increase in the value of the EVPA from 78° to 315° (2008 March 10–2008 July 9) and a large rotation of ~237° during Cycle I, coinciding with the 2008 flare.

The large rotation of EVPA can be interpreted as due to an asymmetric distribution of the magnetic field with respect to the jet axis. Massaro et al. (2001) show that the jet has a spiral structure at 1.6 and 5 GHz. On the other hand, Gabuzda et al. (1994) suggest that the polarization degree and the different values of the EVPA from their VLBI images can be due to shocks propagating along a curved jet, producing an ordered magnetic field with helical structure. These studies found that the jet of W Comae has a projected position angle of ~110° at 1.6 GHz and 5 GHz. Therefore, the rotation can be produced by a swing of the jet along the visual line of sight, or a curved trajectory of the dissipation/emission pattern. In agreement with Abdo et al. (2010b), the second possibility may be due to the propagation of a knot emission that follows a helical path in a
Figure 6. Temporal variability of some physical parameters related to the relativistic jet of W Comae: (from top to bottom) Doppler factor, $\delta(t)$; viewing angle of the jet, $\Phi(t)$; rest-frame viewing angle of the shock, $\Psi(t)$; compression factor of the shocked to the unshocked plasma, $\eta(t)$. These parameters were estimated using the maximum flux value observed during Cycle I. Finally, the lower panel shows the degree of linear polarization, $p\%$, in the R band.

magnetically dominated jet or can be due to an entire bending of the jet.

The direct association found between the $\gamma$-ray flare in 2008 and the gradual change in the EVPA suggests that the $\gamma$-ray and optical emission regions are co-spatial. This implies a highly ordered magnetic field in regions where the $\gamma$-rays emission is produced, therefore this strong flare could have been produced by a strong shock. Taking into account the properties mentioned above, we assumed that the strong flare observed in 2008 in optical and in $\gamma$-rays is a combination of two factors. On one hand, if a curved structure of the jet is assumed, the jet direction will be oriented toward the observer with the minimum viewing angle. On the other hand, a strong shock occurred at the same time. We will discuss this hypothesis in the following section.

4.1. Alignment of Magnetic Field by the Shock

From our results, the moderate anti-correlation found in some flares between the flux and the polarization degree indicates that the magnetic field tends to be aligned with the jet. This result is in agreement with Gabuzda & Cawthorne (1996). However, during the 2008 June major flare, lasting in $\gamma$-rays only three days, the flux correlates with the polarization degree thus suggesting that this event was originated by a transversal shock.

Thus in the observer’s reference frame, the flux of the shocked region is amplified as

$$ F = F_0 \nu^{-\alpha} \delta^{3(\alpha)} , $$

(see Dermer & Menon 2009) where $\delta = [\Gamma_j(1 - \beta \cos \Phi)]^{-1}$ is the jet’s Doppler factor, $\beta = (1 - \Gamma_j^{-2})^{1/2}$ its global velocity in units of speed of light, $\Phi$ is the viewing angle, and $\alpha$ is the spectral index in the optical bands.

The observed degree of polarization $p$ depends on the rest-frame angle between the line of sight and the compression axis $\Psi$, the spectral index $\alpha$, and the shock compression factor $\eta$, which is the ratio of densities of a plasma of relativistic electrons from the shocked to the unshocked region $\eta = \eta_{\text{shock}}/\eta_{\text{unshock}}$ (Hughes & Miller 1991):

$$ p \approx \frac{\alpha + 1}{\alpha + 5/3} \frac{(1 - \eta^{-2}) \sin^2 \Psi}{2 - (1 - \eta^{-2}) \sin^2 \Psi} , $$

(10)

and

$$ \Psi = \tan^{-1} \left( \frac{\sin \Phi}{\Gamma_j(\cos \Phi - \sqrt{1 - \Gamma_j^{-2}})} \right) . $$

(11)

Following Acciari et al. (2009), we assumed a bulk Lorentz factor $\Gamma_j = 20$ for W Comae. We also used the value of $\alpha = 0.87$, given by Tosti et al. (1998).

From Equation (9), we can estimate the Doppler factor as a function of time. The value of $F_0$ is determined by $F_0 = F_{\text{max}} \nu^{\alpha}/\delta_D^{3(\alpha)}$, where $F_{\text{max}}$ is the maximum observed flux and $\delta_D$ is obtained from $\Phi_0$, which is calculated from Equations (10) and (11) for $p \approx p_{\text{var}}$, this being the maximum value of the polarization degree of the variable component (see Table 4). From Hughes & Miller (1991), for $\Psi = \pi/2$, $\eta = 2.2$, which is the minimum compression that produces a degree of linear polarization as high as 45%. This yields to $\Psi_0 \approx 70^\circ$, $\Phi_0 \approx 2.0^\circ$, and $\delta_D \approx 26.7$ at the maximum polarization of the variable component. Using Equations (9), (10), and (11) the
of the Doppler factor can produce large flux variations while 
\[ \theta \approx 65^\circ \] 
Assuming that the position angle of the radio jet is 
\[ \Delta \theta \approx 11^\circ \] 
The viewing angle of the jet \( \Phi \), shows a minimum value of 2:0 and a maximum value of 3:6, i.e., \( \Delta \Phi \approx 1.6^\circ \). These small variations of the Doppler factor can produce large flux variations while \( \Gamma_j \) remains constant. In the state of maximum brightness, the viewing angle of the shock \( \Psi \approx 70^\circ \) undergoes its maximum aberration due to relativistic effects.

A maximum compression of the plasma of \( \eta = 1.69 \) is found, when the polarization degree observed reaches its maximum value of 33\% (2013 May 17 or JD 2456429). The minimum compression factor \( \eta_{\text{min}} = 1.01 \) is obtained when the polarization degree had a minimum value of 0.6\% (2008 July 9 or JD 2454656). These small changes in the compression factor \( (\Delta \eta \approx 0.68) \) can produce large changes in the polarization.

The Doppler factor \( \delta_D \) is obtained when the \( R \)-band flux is at its maximum value due to the presence of the shock. Then, the change in the magnetic field intensity due to the shock is estimated assuming that the minimum variability timescale is related to the shock-front thickness. This scale is estimated considering the lifetime of the synchrotron electrons (see, e.g., Hagen-Thorn et al. 2008). The lifetime of the synchrotron electrons for a given frequency \( \nu \) in GHz is

\[
 t_{\text{loss}} = 4.75 \times 10^2 \left( \frac{1 + z}{\delta_D v_{\text{GHz}} B^3} \right)^{1/2} \text{days},
\]

where \( B \) is the magnetic field in Gauss. Since \( t_{\text{loss}} \approx t_{\text{var}} \), for \( \delta_D = 26.7 \) and \( t_{\text{var}} = 3.3 \pm 0.3 \) days, Equation (12) yields an estimate of the magnetic field intensity, \( B = 0.12 \pm 0.01 \) G, and an upper limit for the emission region size of \( r_{\gamma} \leq c t_{\text{var}} \delta_D / (1+z) = (2.1 \pm 0.2) \times 10^{17} \) cm.

Finally, in cycles where no correlation was found between the flux and the polarization degree, the flares could be possibly due to an oblique shock to the jet’s direction, or due to changes in the Doppler factor, related to changes in the viewing angle of the jet. Therefore, three scenarios are proposed to explain the flares observed at different timescales: (1) a shock transverse to the jet axis, ordering the magnetic field parallel to the shock’s plane; (2) an oblique shock with respect to the jet axis produced in an initially disordered magnetic field, produces a final magnetic field with a component almost parallel to the jet axis; and (3) variations of the Doppler factor due to changes in the jet axis orientation with respect to the observer’s line of sight.

From the polarimetric analysis, we found that the behavior of the polarized variable flux could be due to the superposition of two optically thin synchrotron components. One stable with \( \theta_{\text{cons}} \approx 65^\circ \), \( p_{\text{cons}} \approx 11\% \) and the other variable (see Figure 5). Assuming that the position angle of the radio jet is \( \theta_{\text{jet}} \approx 110^\circ \), we propose that the transversal shocks to the jet axis could be related to the variable component and the oblique shocks to the stable component. Nevertheless, both variable and stable components can be affected by variations of the Doppler factor.

5. CONCLUSIONS

From the photopolarimetric observations of W Comae, we found that the source displayed activity during the monitored period. We clearly detect four flares, estimating that the object has a minimum variability scale of 3.3 days and a maximum variability in brightness of 2.27 mag. The maximum degree of linear polarization reached by W Comae during the campaign was 33.8\%.

An important observational result is the large rotation of EVPA of \( \Delta \theta \approx 237^\circ \), associated with the optical flare and coincident with the major \( \gamma \)-ray flare observed in 2008 June. Subsequently, the polarization angle tends to a preferential orientation of \( \sim 65^\circ \). The large rotation associated with the flare in \( \gamma \)-rays suggests that both optical/\( \gamma \)-ray emissions could be produced in the same jet’s region.

From the analysis of the Stokes parameters, we infer the presence of two optically thin synchrotron components with different polarimetric characteristics: one is a variable component and the other one is stable with a constant degree of polarization of \( p_{\text{cons}} \approx 11\% \) and a constant position angle of \( \theta_{\text{cons}} \approx 65^\circ \). Assuming that the 2008 June optical flare has originated in a transversal shock propagating down a twisted jet and that the source is a spherical blob of radius \( r_{\gamma} \) moving with a Lorentz’s factor of \( \Gamma = 20 \), from our polarimetric data, we estimated a Doppler factor of \( \delta_D \approx 27 \) when the flux was maximum and a visual angle of the jet \( \Phi \approx 2^\circ \). We also obtained a magnetic field intensity \( B \approx 0.12 \) G. Finally, an upper limit for the size of the emission region of \( r_{\gamma} \leq 2 \times 10^{17} \) cm was estimated.

The variability timescales displayed by W Comae show two main characteristics. (1) There are two components in the light-curve, one contributing to the long-term brightening with timescales going from two to four months, and the other that contributes to the short timescale variations (\( \sim 3 \) days). This result is in agreement with Tosti et al. (2002). (2) The Doppler factor changes \( (\delta(t) \approx 16–27) \) could be due to changes in the viewing angle of the jet, implying flux variations lasting \( \sim 0.9 \) yr.

Based on the anticorrelation found between the polarization percentage and the flux, we propose that the observed long-term flux behavior can be explained with a spiral jet and a transversal shock-wave models. This anticorrelation depends on the Doppler factor time variations \( \delta(t) \) for a range of values of the viewing angle \( \theta(t) \).

From our observations, we found that the EVPA in the optical has a value \( \sim 110^\circ \) on 2008 June 7, a day before the \( \gamma \)-ray flare. This value is identical to the projected angle of the radio jet found by Gabuzda et al. (1994) and Massaro et al. (2001). Later, after the strong \( \gamma \)-ray flare finished, the EVPA increases its value up to \( 315^\circ \) on 2008 July 9, and two days afterward it goes down to 229 deg. During the following cycles, the EVPA shows a preferential value of \( 65^\circ \). In a future work, it would be useful to measure the direction of EVPA rotations using also radio data. This will allow us to verify whether the behavior of the EVPA in the optical bands is similar to the EVPA variations studied in the radio bands.

On the other hand, although Zhang et al. (2008) predicts a flare around 2013, we did not detect any important outburst in W Comae from our data collected in 2013. Rather, we report a continuous brightness decrease detected since the beginning of 2008, reaching a minimum value in 2013 May. However, this could also be considered as a prelude to a major flare or a flare that could start at the end of 2013.

We thank the anonymous referee for helpful comments that improved the presentation of this work. M.S., E.B., D.H., J.I.C., and R.M. acknowledge financial support from UNAM–DGAPA–PAPIIT through grant IN116211 and E.B. through grant IN111514. We thank the OAN–SPM staff for...
the support given to this project. This research has made use of the SAO/NASA Astrophysics Data System (ADS) and of the NASA/IPAC Extragalactic Database (NED), which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration.

REFERENCES

Abdo, A. A., Ackermann, M., Agudo, I., et al. 2010a, ApJ, 716, 30
Abdo, A. A., Ackermann, M., Ajello, M., et al. 2010b, Natur, 463, 919
Acciari, V. A., Aliu, E., Aune, T., et al. 2009, ApJ, 707, 612
Acciari, V. A., Aliu, E., Beilicke, M., et al. 2008, ApJL, 684, L73
Agudo, I., Marscher, A., Jorstad, S. G., & Gómez, J. L. 2013, in Proc. of the X
Scientific Meeting of the Spanish Astronomical Society, ed. J. C. Guirao et al. (Barcelona: SEA), 152
Agudo, I., Thum, C., Wiesemeyer, H., & Krichbaum, T. P. 2010, ApJS, 189, 1
Angel, J. R. P., & Stockman, H. S. 1980, ARA&A, 18, 321
Babadzhanyants, M. K., & Belokon’, E. T. 2002, ARep, 46, 609
Belokon’, É. T., Babadzhanyants, M. K., & Pollock, J. T. 2000, A&A, 356, L21
Blandford, R. D., & Rees, M. J. 1978, in Proc. of Pittsburgh Conference on BL
Lac Objects, ed. A. M. Wolfe (Pittsburgh, PA: Univ. of Pittsburgh), 328
Böttcher, M., Mukherjee, R., & Reimer, A. 2002, ApJ, 581, 143
Browne, I. W. A. 1971, Natur, 231, 515
Burbridge, G. R., Jones, T. W., & OdeLL, S. L. 1974, ApJ, 193, 43
Costamante, L. 2012, MmSAI, 83, 138
Dermer, C. D., & Menon, G. 2009, in Princeton Series in Astrophysics, High
Energy Radiation from Black Holes: Gamma Rays, Cosmic Rays, and
Neutrinos, ed. P. U. Pres (Princeton, NJ: Princeton Univ. Press), 62
Fiorucci, M., & Tosti, G. 1996, A&AS, 116, 403
Fossati, G., Celotti, A., Ghisellini, G., & Maraschi, L. 1997, MNRAS, 289, 136
Gabuzda, D. C., & Cawthorne, T. V. 1996, MNRAS, 283, 759
Gabuzda, D. C., Cawthorne, T. V., Roberts, D. H., & Wardle, J. F. C. 1992, ApJ, 388, 40
Gabuzda, D. C., Mullan, C. M., Cawthorne, T. V., Wardle, J. F. C., & Roberts, D. H. 1994, ApJ, 435, 140
Hagen-Thorn, V. A., Larionov, V. M., Jorstad, S. G., et al. 2008, ApJ, 672, 40
Hartman, R. C., Bertsch, D. L., Bloom, S. D., et al. 1999, ApJS, 123, 79
Heidt, J., & Wagner, S. J. 1996, A&A, 305, 42
Holmes, P. A., Brand, P. W. J. L., Impey, C. D., et al. 1984, MNRAS, 211, 497
Hovatta, T., Valtaoja, E., Tornikoski, M., & Lahteenmäki, A. 2009, A&A, 494, 527
Hughes, P. A., & Miller, L. 1991, in Cambridge Astrophysics Series, Vol. 19,
Beams and Jets in Astrophysics, ed. P. A. Hughes (Cambridge: Cambridge
Univ. Press), 1
Jones, T. W., Rudnick, L., Aller, H. D., et al. 1985, ApJ, 290, 627
Jorstad, S. G., Marscher, A. P., Lister, M. L., et al. 2005, AJ, 130, 1418
Liu, F. K., Xie, G. Y., & Bai, J. M. 1995, A&A, 295, 1
Marscher, A. P., Jorstad, S. G., Larionov, V. M., et al. 2010, ApJL, 710, L126
Massaro, E., Maesano, M., Montagni, F., et al. 1999, A&A, 342, L49
Massaro, E., Montovani, F., Fantò, R., et al. 2001, A&A, 374, 435
Nieppola, E., Tornikoski, M., & Valtaoja, E. 2006, A&A, 445, 441
Nilsson, K., Pasanen, M., Takalo, L. O., et al. 2007, A&A, 475, 199
Nilsson, K., Pursimo, T., Heidt, J., et al. 2003, A&A, 400, 95
Padovani, P., & Giommi, P. 1995, ApJ, 444, 567
Sasada, M., Uemura, M., Fukazawa, Y., et al. 2011, PASJ, 63, 489
Schmidt, G. D., Elston, R., & Lupie, O. L. 1992, AJ, 104, 1563
Sorcia, M., Benítez, E., Hiriart, D., et al. 2013, ApJS, 206, 11
Tagliaferri, G., Ghisellini, G., Giommi, P., et al. 2000, A&A, 354, 431
Tosti, G., Fiorucci, M., Luciani, M., et al. 1998, A&AS, 130, 109
Tosti, G., Massaro, E., Nesci, R., et al. 2002, A&A, 395, 11
Zhang, X., Zheng, Y. G., Zhang, L., et al. 2008, PASJ, 60, 145