The origin of the local 1/4–keV X-ray flux in both charge exchange and a hot bubble

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The solar neighbourhood is the closest and most easily studied sample of the Galactic interstellar medium, an understanding of which is essential for models of star formation and galaxy evolution. Observations of an unexpectedly intense diffuse flux of easily absorbed 1/4-kiloelectronvolt X-rays1–4, coupled with the discovery that interstellar space within about a hundred parsecs of the Sun is almost completely devoid of cool absorbing gas3, led to a picture of a ‘local cavity’ filled with X-ray-emitting hot gas, dubbed the local hot bubble5–7. This model was recently challenged by suggestions that the emission could instead be readily produced within the Solar System by heavy solar-wind ions exchanging electrons with neutral H and He in interplanetary space8–10, potentially removing the major piece of evidence for the local existence of million-degree gas within the Galactic disk11–13. Here we report observations showing that the total solar-wind charge-exchange contribution is approximately 40 per cent of the 1/4-keV flux in the Galactic plane. The fact that the measured flux is not dominated by charge exchange supports the notion of a million-degree hot bubble extending about a hundred parsecs from the Sun.

When the highly ionized solar wind interacts with neutral gas, an electron may hop from a neutral to an outer orbital of an ion, in what is known as charge exchange. The electron then cascades to the ground state of the ion, often emitting soft X-rays in the process16. The calculations of X-ray intensity from solar-wind charge exchange depend on limited information about heavy ion fluxes and even more uncertain atomic cross-sections. The ‘Diffuse X-rays from the Local galaxy’ (DXL) sounding rocket mission17 was launched from the White Sands Missile Range in New Mexico, USA, on 12 December 2012 to make an empirical measurement of the charge exchange flux by observing a region of higher interplanetary neutral density (with a correspondingly higher charge exchange rate) called the ‘helium focusing cone’. Neutral interstellar gas flows at about 25 km s−1 through the Solar System owing to the motion of the Sun through a small interstellar cloud. This material, mostly hydrogen atoms but about 15% helium, flows from the Galactic direction (longitude l, latitude b) = (3°, 16°), placing Earth downstream of the Sun in early December18. The trajectories of the neutral interstellar helium atoms are governed primarily by gravity, executing hyperbolic Keplerian orbits and forming a relatively high-density focusing cone downstream of the Sun about 6° below the ecliptic plane (Fig. 1)19. Interstellar hydrogen, on the other hand, is also strongly affected by radiation pressure and photoionization: radiation pressure balances gravity, reducing the focusing effect, while photoionization creates a neutral hydrogen cavity around the Sun.

The early December launch of DXL placed the He focusing cone near the zenith at midnight. The 7° field of view was scanned slowly back and forth across one side of the cone and more rapidly in a full circle to test the consistency of the derived charge exchange contribution in other directions and to make measurements of the detector particle background while DXL was looking towards Earth (Extended Data Fig. 1). Figure 2 shows the ROSAT All Sky Survey 1/4-keV map10 (R12 band) with the paths of the DXL slow scan (red) and fast scan (white) overplotted. The ROSAT observation of the slow-scan region was performed with the slow portion shown in red. The black line is the 90° horizon for the DXL flight. The width of the band represents the half-power diameter of the instrument beam.
in September 1990 when the line of sight was about one astronomical unit (the Earth–Sun distance, 1 AU) away from, and parallel to, the He cone, so its charge exchange contribution was not strongly affected by the cone enhancement (Fig. 1).

For this analysis, we chose pulse height limits for both of the DXL proportional counters (Counter-I and Counter-II) to match the pulse heights of the ROSAT 1/4-keV band as closely as possible (Extended Data Fig. 2). This energy range is dominated by and contains most of the emission from solar-wind charge exchange and/or the local hot bubble. To quantify the solar-wind charge exchange emission we compared both DXL and ROSAT count rates with well determined models of the interplanetary neutral distribution along the lines of sight for both sets of measurements (Fig. 3). Figure 3 shows the combined Counter-I and Counter-II count rates (black dots) during the DXL scan and the ROSAT 1/4-keV count rates in the same directions (blue solid line). The best fit to the DXL total count rate (red solid line), and the solar-wind charge exchange contributions to DXL (red dashed line) and ROSAT (blue dashed line) are also shown (see Table 1 for best-fit parameters: the model shown corresponds to the second column). There is potentially an additional contribution from charge exchange between the solar-wind ions and the geocoronal hydrogen surrounding Earth, which tracks the short-term variations in solar-wind flux. Time variations of a few days or less were removed from the ROSAT maps, and the current best estimate of the residual from geocoronal charge exchange is about 50 ROSAT units (1 RU = 10⁻⁶ counts s⁻¹ arcmin⁻²) for the ROSAT 1/4-keV band (K.D.K., J. Carter, M.P.C., Y. M. Colladovega, M.R.C., T.E.C., D.K., F.S.P., A. Read, I.P.R., D. G. Sibeck, S. F. Sembay, S.L.S., N.E.T. and D.M.W., manuscript in preparation). The geocoronal contribution to the DXL flux should be negligible, owing to the look direction, which is directly away from the Sun. The signature of the cone enhancement in the DXL data compared to the ROSAT rates is evident, highlighting the contribution from charge exchange. However, the best fit shows that the total charge exchange contribution to ROSAT is only about 40% ± 5% (statistical error) ± 5% (systematic error) of the total flux observed at the Galactic plane. Its contribution to the ROSAT flux over the DXL scan path is typically about 140 RU. For comparison, the total ROSAT 1/4-keV flux ranges from around 300 RU to 400 RU in the plane up to 1,400 RU in the brightest areas at intermediate and high latitudes. This result implies that the measured fluxes are dominated by interstellar emission, strengthening the original idea of a hot bubble filling the local interstellar medium for a hundred parsecs or so in all directions from the Sun.

It has been pointed out that a hot bubble creates an apparent pressure balance problem with the tenuous warm cloud that the Sun is passing through. However, recent results on the magnetic contribution to the cloud pressure and new three-dimensional maps of the local interstellar medium bring the implied pressure of the plasma in the local hot bubble to rough agreement with pressures derived for the local interstellar clouds when the measured contribution from the solar-wind charge exchange is removed from the local hot bubble emission.

Online Content Methods, along with any additional Extended Data display items and Source Data, are available in the online version of the paper; references unique to these sections appear only in the online paper.

Received 20 March; accepted 14 May 2014. Published online 27 July 2014.

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**Figure 3** Neutral column density for DXL and ROSAT. Neutral column density distribution integrals for each line of sight along the scan path. The density distribution in the integrals is weighted by one over the distance from the Sun squared (1/R²) to reflect the dilution of the solar wind as it flows outward. The red lines are the integral for He (solid) and H (dashed) in the DXL geometry. The blue lines represent the integral for He (solid) and H (dashed) in the ROSAT geometry. The black line shows the Galactic latitude during the geometry. The blue lines represent the integral for He (solid) and H (dashed) in the DXL geometry. The black line shows the Galactic latitude during the geometry. The blue lines represent the integral for He (solid) and H (dashed) in the ROSAT geometry. The black line shows the Galactic latitude during the geometry. The blue lines represent the integral for He (solid) and H (dashed) in the ROSAT geometry. The black line shows the Galactic latitude during the geometry.

**Figure 4** Fit to DXL and ROSAT data. Combined Counter-I and Counter-II count rates (black dots) during the DXL scan and ROSAT 1/4-keV count rate in the same directions (blue solid line). The best fit to the DXL total count rate (red solid line), and the solar-wind charge exchange contribution to DXL (red dashed line) and ROSAT 1/4-keV bands (blue dashed line) are also shown. The error bars are s.e.m.

**Table 1** Best-fit model parameters

| Parameter | ROSAT geocoronal contribution (RU) | Total solar-wind charge exchange contribution to ROSAT R12 data |
|-----------|-----------------------------------|---------------------------------------------------------------|
| n₂H/R₂_Ha | 1                                | 0.63 ± 0.13                                                  |
| Counter-I | 0.91 ± 0.06                       | 42 ± 6%                                                      |
| Counter-II| 0.97 ± 0.02                       | 6%                                                          |

Summary of best-fit parameters for different assumptions for the geocoronal contribution to the ROSAT R12 band, and the ratio between hydrogen and helium compound cross-sections n₂H/a_R₂. Counter-I and Counter-II are the ratios of the fitted DXL response to the nominal value from laboratory calibrations (the corrections are well within the range expected from spectral uncertainties). u_rel is the relative speed between solar wind and neutral flow and n₂H is the proton density. For the ratio of solar-wind fluxes during the DXL and ROSAT measurements, we note that although both missions were near the solar maximum (and therefore should have similar isotropic composition), solar activity in terms of sunspots was weaker during the DXL measurement, as reflected by the fitted ratios. The total solar-wind charge exchange contribution to ROSAT (interplanetary + geocoronal) is divided by the observed R12 rate and given as a percentage of R12 at b = 0° (333 RU, which is the lowest anywhere on the scan and close to the lowest on sky). Errors are 1σ.
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Acknowledgements We thank the personnel at NASA’s Wallops Flight Facility and the White Sands Military Range for their support of payload development, integration and launch, and the technical personnel at the University of Miami, NASA’s Goddard Space Flight Center and the University of Michigan for their support of the instrument’s development. This work was supported by NASA award numbers NNX11AF04G and NNX09AF09G. D.K. and R.L. acknowledge financial support for their activity through the programme ‘Soleil Héliosphère Magnétosphère’ of the French space agency CNES, and the National Program ‘Physique Chimie du Milieu Interstellaire’ of the Institut National des Sciences de l’Univers (INSU). M.C. and N.E.T. are employed through the Center for Research and Exploration in Space Science and Technology (CRESST) and the University of Maryland, Baltimore County, Baltimore, Maryland, USA.

Author Contributions Y.U., N.E.T., M.G., D.M., M.R.C., F.S.P. and S.T.L. contributed to hardware development. Y.U., N.E.T., M.G., M.R.C., F.S.P., M.C., D.K. and K.M. contributed to launch operations. M.G., D.M., Y.U., N.E.T., M.R.C., D.K., K.D.K., K.M., T.C., I.P.R., S.L.S., E.U. and B.M.W. contributed to data reduction and analysis. D.K. and R.L. prepared the neutral integral distributions. All authors discussed the results and commented on the manuscript.

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METHODS

The total count rate due to charge exchange with H and He is the integral along the line of sight of the product of solar-wind ion flux \( n_i v_{rel} \) and the donor densities \( n_H \) and \( n_{He} \), as functions of position, and the sum over the products of the partial cross-sections for producing each X-ray line by charge exchange and the efficiency of producing a detector count from that line:

\[
\text{Rate} = \sum_i \sum_j n_i n_{He} v_{rel} \sigma_{i,j} b_j g_j ds + \sum_i \sum_j n_i n_{He} v_{rel} \sigma_{i,j} b_j g_j ds
\]

where \( i \) represents the solar-wind species, \( j \) the emission lines for that species, \( \sigma \) are the speed-dependent interaction cross-sections for individual species, \( b_j \) is the line branching ratio, \( g_j \) is the instrument’s efficiency for detecting photons in line \( j \), and \( v_{rel} \) is the relative speed between solar wind and neutral flow (both the bulk and thermal velocity). We can then write the ion density \( n_i \) in terms of the proton density \( n_p \), at \( R_0 = 1 \) AU, assuming that it scales as one over the square of the distance \( R \) from the Sun (we verified that neutralization effects on the solar-wind ions are negligible), and define the compound cross-section as:

\[
\alpha = \sum_i \sum_j \frac{n_i(R_0)}{n_p(R_0)} \sigma_{i,j} b_j g_j
\]

In the case of ‘constant’ solar-wind conditions, the solar-wind flux can be removed from the integrals, and the total charge exchange rate with H and He can be written as:

\[
\text{Rate} = n_p(R_0) v_{rel} \left( \frac{\alpha_i}{R^2} \frac{\text{d}n}{\text{d}t} + \frac{\alpha_i}{R^2} \frac{\text{d}n}{\text{d}t} \right)
\]

The assumption of isotropy of the ion flux included in the equation above is an approximation, because the flux is known to vary strongly on timescales of about a day. Evidence that averaging over the few-week transit time through the relevant interplanetary region smooths these fluctuations is found in the very good agreement between four complete sky surveys performed years apart by different missions. A factor to account for the difference in the solar-wind flux between the ROSAT and DXL is included in our fitting procedure and its best value is reported in Table 1.

In this work, we used the expected H and He distribution adapted for solar-wind conditions during the respective missions to calculate the integrals for both H and He for each point along the DXL scan path. The distribution of the interplanetary neutrals is calculated based on the solar ionization conditions derived from measurements of backscattered solar radiation and checked by in situ sampling, so the integrals above can be calculated with some confidence for all lines of sight. We took as free parameters the combination \( n_p(R_0) v_{rel} \alpha_{He} \), the ratio of solar-wind fluxes for the two missions \( (n_p(R_0) v_{rel} \alpha_{DXL})/(n_p(R_0) v_{rel} \alpha_{ROSAT}) \), and correction factors to fine-tune the calculated ratios of DXL counter responses to the ROSAT 1/4-keV band response. We then did a global least-squares fit for both DXL counter rates for each point along the scan path.

There is insufficient variation in the hydrogen column densities to determine its effective cross section, so for \( \sigma_{He} \), we assumed ratios of one, as assumed in the calculated contributions, and two, since some determinations show smaller cross-sections for helium. The residual contribution to ROSAT from charge exchange in Earth’s geocorona has in the past been assumed to be negligible owing to the good agreement with other all-sky surveys done with different observing geometries. More recent analyses suggest that the value should be more like 50 RU. Table 1 shows the best-fit parameters for some different assumptions for the residual geocoronal charge exchange contribution to the ROSAT rates and for the ratio of effective cross-sections for hydrogen and helium. The total solar-wind charge exchange contribution is minimally affected. A systematic error has been included in our results to account for the variation in the table.

There may be other potential sources of systematic uncertainties affecting our result. These include contribution from point sources to the DXL count rate, scattered solar X-rays, higher-energy photons into the DXL bands, and a highly non-isotropic solar-wind flux. We estimated the first three and found their contribution to our result to be within a few per cent; for the fourth, there is no evidence of a non-isotropic solar wind across the DXL slow scan that would significantly (within error) affect our result.
*Extended Data Figure 1* | DXL count rates versus time during flight. Count rate of Counter-I (red) and Counter-II (black) as a function of time during launch. The DXL observation started off the cone, moved towards the nose of the cone (first scan) and back (second scan) at about 0.7 degrees per second, then performed an Earth scan at about 10 degrees per second to measure the instrument background (fast scan), returned to the nose of the cone, and then performed another scan off the cone (third scan) and back to the nose (fourth scan), for a total of four slow scans along the cone and one fast scan. The gradient up and down the He focusing cone is evident. The error bars are s.e.m.
Extended Data Figure 2 | The DXL detector response. Effective grasp (solid angle times area) for the DXL Counter-I (red) and Counter-II (black), compared with the ROSAT 1/4-keV (R12 band) effective area (blue).