LOW MILKY-WAY-LIKE MOLECULAR GAS EXCITATION OF MASSIVE DISK GALAXIES AT $z \sim 1.5$

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ABSTRACT

We present evidence for Milky-Way-like, low-excitation molecular gas reservoirs in near-IR-selected massive galaxies at $z \sim 1.5$, based on IRAM Plateau de Bure Interferometer CO[3–2] and NRAO Very Large Array CO[1–0] line observations of two galaxies that had been previously detected in CO[2–1] emission. The CO[3–2] flux of BzK $-21000$ at $z = 1.522$ is comparable within the errors to its CO[2–1] flux, implying that the CO[3–2] transition is significantly subthermally excited. The combined CO[1–0] observations of the two sources result in a detection at the $3\sigma$ level that is consistent with a higher CO[1–0] luminosity than that of CO[2–1]. Contrary to what is observed in submillimeter galaxies and quasi-stellar objects, in which the CO transitions are thermally excited up to $J \geq 3$, these galaxies have low-excitation molecular gas, similar to that in the Milky Way and local spirals. This is the first time that such conditions have been observed at high redshift. A Large Velocity Gradient analysis suggests that molecular clouds with density and kinetic temperature comparable to local spirals can reproduce our observations. The similarity in the CO excitation properties suggests that a high, Milky-Way-like, CO to H$_2$ conversion factor could be appropriate for these systems. If such low-excitation properties are representative of ordinary galaxies at high redshift, centimeter telescopes such as the Expanded Very Large Array and the longest wavelength Atacama Large Millimeter Array bands will be the best tools for studying the molecular gas content in these systems through the observations of CO emission lines.

Key words: cosmology: observations – galaxies: evolution – galaxies: formation – galaxies: high-redshift – galaxies: starburst

1. INTRODUCTION

Daddi et al. (2008; D08 henceforth) recently reported the discovery of large molecular gas reservoirs inside near-IR-selected galaxies with star formation rates SFR $\sim 100–150$ $M_\odot$ yr$^{-1}$ and stellar masses $M^* \sim 10^{11}$ $M_\odot$ at redshift $z \sim 1.5$. The SFR in these systems is an order of magnitude lower than in submillimeter galaxies (SMGs) and quasi-stellar objects (QSOs) that also were detected previously in molecular gas emission at millimeter galaxies (SMGs) and quasi-stellar objects (QSOs) in the distant universe have very dense and warm gas that is typical spirals (Solomon & Vanden Bout 2005). Despite many objects are similar to those of local starbursts and ULIRGs. Given that both galaxies observed in CO by D08 were detected with large CO luminosities, and that the space density of similar sources is $10–30$ times larger than that of SMGs (Daddi et al. 2007), this suggests that such a quiescent gas-consumption activity could be a prevalent mode of star formation and galaxy growth in the distant universe.

In order to gain insight to the physical properties of the molecular gas and thus on the nature of the star formation in these galaxies it is important to investigate the CO excitation properties. In galaxies like the Milky Way (Fixsen et al. 1999) and local star-forming spiral galaxies (e.g., Braine & Combes 1992; Young et al. 1995; Mauersberger et al. 1999; Crosthwaite & Turner 2007; Papadopoulos & Seaquist 1998) the molecular gas has relatively low excitation and is rather diffusely distributed (low density and temperature), resulting in a high CO to H$_2$ conversion factor. In contrast, dusty starburst systems like ULIRGs in the local universe and SMGs and QSOs in the distant universe have very dense and warm gas (Tacconi et al. 2006, 2008; Weiß et al. 2005a, 2005b; Riechers et al. 2006; Downes & Solomon 1998), with a CO spectral energy distribution showing thermalized emission up to at least the rotational transition $J = 3$ (Weiß et al. 2007b), and a CO to H$_2$ conversion factor of about five times higher than that of local galaxies.

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study, and by the prevalence of high frequency, high-J transition observations.

Here, we present multi-J CO observations based on PdBI and VLA observations for two $z \sim 1.5$ BzK-selected galaxies previously detected in CO[2–1] with the aim of estimating the excitation properties of CO in these sources.

2. OBSERVATIONS

2.1. Plateau de Bure Observations

We observed BzK − 21000 at $z = 1.522$ in the CO[3–2] transition (rest-frame 345.796 GHz) with the PdBI in D-configuration in 2008 May and June. The on-source observing time was 9.7 hr with the full array (six antennae) and a beam size of $4''0 \times 3''.1$. We tuned the receivers to 136.970 GHz (2 mm window), slightly offset to the frequency of the CO[3–2] line at $z = 1.522$ at 137.112 GHz. The total bandwidth of our dual polarization mode observations was 1 GHz (corresponding to $\sim 2200$ km s$^{-1}$). The correlator has eight independent units each covering 320 MHz (128 channels each with a width of 2.5 MHz). The phase center was 7.5 away from the position of BzK − 21000 in order to target the SMG GN20 during the same observations (Daddi et al. 2009). All the values quoted below were corrected for a primary beam attenuation (PBA) of 11%. We reduced and calibrated the data with the GILDAS software packages CLIC and MAP. The final data, imaged using natural weighting, have a noise of 0.11 mJy integrated over the full 1 GHz bandwidth. In addition to the CO[2–1] data set already presented by D08, we also use CO[2–1] observations of both BzK galaxies obtained with IRAM PdBI in B-configuration (synthesized beam $\sim 1''.3$) that have been already presented in Daddi et al. (2009). The new data solidly confirm the CO[2–1] detections of D08. Our flux calibration in both the 2 mm and 3 mm observations with IRAM PdBI is primarily based on measurements of the standard calibrator MWC349 for which accurate models of the spectrum at these frequencies are available. The knowledge of the typical ranges of PdBI antennas efficiencies at the various frequencies can also be used as a cross check of the flux calibration. Overall, we estimate that the typical accuracy on the absolute flux calibration for our observations are within a range of $\sim 10\%$ in the 3 mm band and $\sim 22\%$ in the 2 mm band. To the measurement errors, we add in quadrature uncertainties corresponding to half those ranges for the estimate of the total error budget on fluxes and luminosities in each observing transition.

2.2. Very Large Array Observations

We observed both sources—BzK − 21000 and BzK − 4171—with the VLA in D-configuration in 2008 August and September (synthesized beam of $\sim 1''.5$). In order to measure the CO[1–0] transition (rest-frame 115.271 GHz) we used two contiguous 50 MHz IF bands tuned at 45.6851 GHz and 45.7351 GHz for BzK − 21000 and at 46.7351 GHz and 46.7851 GHz for BzK − 4171. The total bandwidth of 100 MHz corresponds to about 650 km s$^{-1}$ in velocity space, which covers well most of the emission lines. The total integration time was 21 hr and 28 hr on-source for BzK − 21000 and BzK − 4171, respectively. We reduced and calibrated the data with standard techniques using AIPS. Given that summer time observations with the VLA can be subject to bad phase stability, we employed fast switching phase calibration using the calibrator J1302 + 578. We also included cycles where we switched to a position near to the phase calibrator, and used these data to check the phase coherence time. From this we measured the flux correction factor due to de-coherence of 24% for BzK − 21000 and 15% for BzK − 4171. We correct our measurements by these factors to account for these small losses of signal. The final maps (100 MHz bandwidth) have rms noise levels of 0.11 mJy and 0.13 mJy for BzK − 21000 and BzK − 4171, respectively, once corrected for de-coherence and PBA at the 10% level. For the VLA observations, by monitoring the bootstrapped flux of the phase calibrator over the observing days, we estimate a gain calibration uncertainty in a range of 14%, small compared to the typical measurement uncertainties for our objects (see the next section).

3. RESULTS

3.1. CO[3–2] in BzK − 21000

The CO[3–2] transition in BzK − 21000 is securely detected in the IRAM PdBI data at the $\sim 6\sigma$ level (Figure 1). In order to properly compare the fluxes from the different CO transitions, and to derive a CO spectral line energy distribution (SLED), we need to perform flux measurements over the same velocity range in the different transitions. We use a velocity range of 525 km s$^{-1}$ centered on the CO[2–1] line (the transition detected with the highest signal-to-noise ratio ($S/N$), Figure 2). We also corrected line fluxes for possible underlying continuum emission. Direct measurements in the PdBI maps return no significant evidence for continuum emission (Figures 1 and 2) with formal measurements of $S(3 \text{ mm}) = 40 \pm 60 \mu\text{Jy}$ and $S(2 \text{ mm}) = 320 \pm 150 \mu\text{Jy}$. Given the estimate of $L_{\text{IR}} \sim 10^{12} L_{\odot}$ in these galaxies (D08) we would expect continuum emission at the level of $S(3 \text{ mm}) = 30 \mu\text{Jy}$ and $S(2 \text{ mm}) = 100 \mu\text{Jy}$ based on the SED models of Chary & Elbaz (2001) that have been shown to accurately reproduce the properties of BzK galaxies (e.g., Daddi et al. 2005, 2007). These predictions are consistent with our measurements. Given the nondetections of the continuum in our observations, we used the model estimates as our best guesses and subtracted those values to obtain continuum-free line measurements. These continuum corrections are small, 7% and 2.5% for the 2 mm and 3 mm band, respectively. We did not correct the VLA measurements for continuum emission, which is expected to be entirely negligible at 45 GHz ($\ll 10 \mu\text{Jy}$).

Applying all the corrections and adding all sources of uncertainty in quadrature, our CO measurements for BzK − 21000 are $I_{\text{CO(3–2)}} = 0.70 \pm 0.15 \text{ Jy km s}^{-1}$ and $I_{\text{CO(2–1)}} = 0.62 \pm 0.07 \text{ Jy km s}^{-1}$. This new CO[2–1] measurement is slightly lower than that reported in D08 but it has a higher S/N. The CO[2–1] spectrum now clearly shows a double peak profile, similar to BzK − 4171 in D08, suggestive of rotation.

In the following, we express the line ratios in terms of brightness temperature ratios (or equivalently, luminosity ratios) of the different transitions, as can be derived from Equation (3) of Solomon & Vanden Bout (2005): $r_{J-J-1} = I_{\text{CO}(J-1-J-J-1)} / I_{\text{CO}(J-1-J-J-2)} \times ((J-1)^2 / J^2)$. In the case of thermalized transitions the temperature ratio is 1 by definition, the CO luminosity is the same in the different transitions and the CO fluxes in Jy km s$^{-1}$ scale proportionally to $J^2$.

For BzK − 21000, we derive a solid measurement of $r_{32} = 0.50 \pm 0.12$. The $r_{32}$ measurement clearly shows that the CO[3–2] transition is subthermalized with respect to CO[2–1], a result significant at the $> 4\sigma$ level. From Figure 2, it seems that CO[3–2] might be particularly weak in the reddest half of the spectrum, a possible sign of differential excitation.
Figure 1. Top panel: CO[3–2] map centered on BzK − 21000, averaged over 525 km s\(^{-1}\), corresponding to the observed range of CO emission (see also Figure 2). Contour levels start at ±2σ and are in steps of 1σ ([≈ 0.215 mJy beam\(^{-1}\)] with positive (negative) contours shown as solid (dashed) lines. The white cross shows the VLA 1.4 GHz position. Bottom panel: map averaged over the remaining channels.

Measuring \( r_{32} \) for the red and blue halves separately we do find a suggestion for different excitation (\( r_{32} = 0.28 \) and 0.64, respectively), but the difference is significant only at the 1.7σ level.

3.2. CO[1–0] Constraints

Inspection of the VLA observations of the two BzK galaxies shows \( \gtrsim 2\sigma \) peak signal close to the positions of both galaxies. Individual fluxes cannot be reliably derived from the data. In order to increase the S/N in the measurements we corrected each VLA map for de-coherence and PBA, aligned the maps at the VLA 1.4 GHz continuum coordinates of the two galaxies, and averaged the data from the two objects. Figure 3 shows that some positive CO[1–0] signal is detected at the expected position, significant at about 3.1σ and corresponding to a peak flux density of 0.28 ± 0.09 mJy. Given the insufficient S/N of the data, it is impossible to know if the CO[1–0] emission is slightly extended. In order to compare this CO[1–0] measurement with the CO[2–1] results we have used the B-configuration CO[2–1] observations with IRAM PdBI, which have comparable angular resolution to the VLA CO[1–0] data, and extracted the signal over velocity ranges matching the VLA observations on each galaxy (covering both lines very well). This results on an average peak CO[2–1] signal in the combined CO[2–1] map of 0.67 ± 0.09 mJy. We derive \( r_{21} = 0.60^{+0.40}_{-0.17} \), averaged over the two BzK galaxies. Taken at face value, this implies that we are seeing excess peak emission in CO[1–0] relative to what is expected from CO[2–1] although the evidence for an excess is very mild given the low S/N in the CO[1–0] data. Higher S/N observations would be needed to establish if the CO[2–1] transition is subthermally excited with respect to CO[1–0].
LVG models (model 1, 2, 3, respectively). The Milky Way SLED from Fixsen et al. (triangles, 1999) is shown for comparison. The constant brightness temperature transition.

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\[ \text{BzK} = 21000 \]

[Figure 4. Left panel: the CO SLED of BzK = 21000 at z = 1.522 obtained from PdBI and VLA observations. The black solid, dotted, and dashed lines show our LVG models (model 1, 2, 3, respectively). The Milky Way SLED from Fixsen et al. (triangles, 1999) is shown for comparison. The constant brightness temperature relation (red dashed line) is normalized to our CO[2–1] measurement. Right panel: the same as left panel but shown in luminosity units, normalized to the CO[2–1] transition.} \]

4. DISCUSSION

For BzK = 21000, the CO[3–2] and [2–1] fluxes are comparable within the errors. The CO SLED of BzK = 21000 is consistent with the Milky Way disk measurements by Fixsen et al. (1999) (Figure 4), and similar to what is observed in local spirals (e.g., Braine & Combes 1992; Young et al. 1995; Mauersberger et al. 1999; Crosthwaite & Turner 2007). The CO SLED of BzK = 21000 is in stark contrast to what has been generally seen so far in high-redshift sources like QSOs and SMGs, for which gas has been found to be thermalized up to at least the J = 3 transition (Weiß et al. 2005a, 2007a, 2007b; Riechers et al. 2006). A possible exception is the extremely red galaxy HR 10 at z = 1.44 according to the results of Andreani et al. (2000) and Papadopoulos & Ivison (2002) who find comparable fluxes of CO[5–4] and CO[2–1]. However, we note that the excitation in our sources is still lower than that in HR10.

We now discuss the implication of our excitation measurements on the physical properties of the gas, including the kinetic temperature \( T_{\text{kin}} \) and density \( n \) of the H\(_2\) gas in the molecular clouds, by using Large Velocity Gradient (LVG) modeling (Goldreich & Kwan 1974; Scoville & Solomon 1974). We use the collision rates from Flower (2001) with a fixed CO abundance to velocity gradient ratio of \( [\text{CO}] / (dv/dr) = 1 \times 10^{-5}\text{pc} (\text{km s}^{-1})^{-1} \) (Weiß et al. 2005b), an H\(_2\) ortho:para ratio of 3:1 and the cosmic background temperature at \( z = 1.522 \) of 6.89 K.\(^9\) Due to the limited amount of information available, we will not be able to resolve the ambiguity between all the free parameters. For example, choosing a different value of \([\text{CO}] / (dv/dr)\) would also affect the SLED, and would require re-adjusting \( T_{\text{kin}} \) and/or \( n \) to reproduce observed ratios (see, e.g., Weiß et al. 2005b, for a more detailed discussion of the parameter space degeneracies). We present here three representative models, spanning the range of parameter degeneracies in \( T_{\text{kin}} \) and \( n \) (see Figure 4) that reproduce the observed CO SLED of BzK = 21000. Model 1 has a kinetic gas temperature of \( T_{\text{kin}} = 25\) K and molecular cloud density of \( n = 1300\) cm\(^{-3}\), very similar to the molecular gas excitation conditions of the Milky Way. The filling factor of clouds, assuming a disk radius \( r_0 \approx 5\) kpc from our ACS observations (D08), is 2% for this model. Model 2 can still fit the data assuming \( T_{\text{kin}} = 90\) K and \( n = 600\) cm\(^{-3}\) with a similar filling factor to model 1. Model 3 results in a higher filling factor of 8% with a low \( T_{\text{kin}} = 10\) K but a higher \( n = 2500\) cm\(^{-3}\). Typical gas/dust temperatures in nearby spirals/starbursts are 10–30 K (e.g., Braine & Combes 1992; Mauersberger et al. 1999), sometimes up to 40–50 K in the central regions of nuclear starbursts (Downes & Solomon 1998). Normal galactic giant molecular clouds have typical average densities of \( 10^2–10^3\) cm\(^{-3}\) (Scoville et al. 1974; Evans 1999). These values are broadly similar to what is allowed from LVG modeling of our observations. Clearly, we can rule out solutions where the gas is at high density and/or temperature, similar to what is generally found inside local ULIRGs, distant SMGs, and QSOs.

Despite the degeneracies in the inferred physical parameters, the LVG models presented above consistently predict \( r_{21} \approx 0.85 \), higher than our average estimate for the BzK galaxies but consistent within the uncertainties. Robustly confirming an excess CO[1–0] emission over what is predicted by the LVG models shown in Figure 4 would imply the discovery of very cold molecular gas unseen even in the CO[2–1] transition. For the moment, given our limited S/N at the measured \( r_{21} \), we assume that \( r_{21} \approx 0.85 \) and \( r_{32} \approx 0.50 \) are the typical values for our sources. This implies that the CO luminosities derived for our galaxies in D08 from CO[2–1] do not need substantial revisions, exception made for the small 16% increase required when converting CO[2–1] to CO[1–0], as suggested by our LVG modeling. However, our results strongly indicate that when using transitions higher than CO[2–1] only a fraction of the CO[1–0] luminosity is recovered (Figure 4, right), implying that such measurements are not reliable for accurately estimating the total gas mass and SFEs in BzK-like galaxies. This evidence could explain the discrepancy between the observations of Hatsukade et al. (2009) and Tacconi et al. (2008) who failed to detect CO[3–2] in the four optically and UV-selected galaxies at \( 1.4 < z < 2.5 \), while so far we have a high detection rate in CO[2–1] at redshift \( z \approx 1.5 \). A large correction factor of 2.4 is required, in fact, at CO[3–2] in order to recover the total molecular gas mass, based on our results from BzK = 21000.

Given that these two high-redshift galaxies have SFEs and CO excitation properties similar to those in local spirals and in the Milky Way, it is plausible that they may also have a Milky-Way-like molecular conversion factor \( \alpha_{\text{CO}} \) as well. This would imply molecular gas reservoirs of \( \sim 10^{11} M_\odot \) and gas mass fractions \( f_{\text{gas}} \gtrsim 0.6 \). To summarize, our sources provide clear examples of the long-sought after population of low-excitation CO emitters at high-z. Observations of a larger sample of near-IR-selected...
galaxies are required in order to confirm our results in a statistical way. In particular, high S/N observations of CO[1–0] and high-J CO are required to search for additional components of colder and warmer gas, respectively. Much higher S/N data will also be required to search for excitation variations inside the disks, as tentatively suggested by our data and as could be expected in the case of very clumpy gas distribution (e.g., Bournaud et al. 2008).

Given that these are the first excitation measurement inside normal, near-IR-selected disk-like galaxies in the distant universe, it is plausible to speculate that such low-excitation properties might be typical in distant galaxies. ALMA will be a powerful device to study CO[2–1] in $z < 3$ galaxies, but at much higher redshift, the higher order transitions may not be excited in the average galaxy (see also Papadopoulos & Ivison 2002). In those cases, centimeter telescopes, such as the EVLA or the Square Kilometer Array (SKA), will become the primary tools for the study of molecular gas in the earliest, normal galaxies. In distant galaxies with massive amounts of cold, diffuse, low-excitation gas we will be able to detect carbon lines with ALMA (e.g., [C ii] $\lambda 158 \mu m$ redshifted to $> 400 \mu m$ in $z > 1.5$ galaxies). This synergistic combination of observations will offer a powerful tool for interpreting the gaseous content of ordinary galaxies at high redshift.

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