How Are Mixed-Phase Clouds Mixed?

Alexei Korolev1 and Jason Milbrandt1

1Environment and Climate Change Canada, Atmospheric Science and Technology Directorate, Toronto, ON, Canada

Abstract

Mixed-phase clouds are recognized as significant contributors to the modulation of precipitation and radiation transfer on both regional and global scales. This study is focused on the analysis of spatial inhomogeneity of mixed-phase clouds based on an extended data set obtained from airborne in situ observations. The lengths of continuous segments of ice, liquid, and mixed-phase clouds present a cascade of scales varying from $10^3$ km down to a minimum scale of 100 m determined by the spatial resolution of measurements. It was found that the phase composition of mixed-phase clouds is highly intermittent, and the frequency of occurrence of ice, liquid, and mixed-phase regions increases with the decrease of their spatial scales. The distributions of spatial scales have well-distinguished power-law dependencies. The results obtained yield insight into the morphology of mixed-phase clouds and have important implications for improvement in representing subgrid inhomogeneity of mixed-phase clouds in weather and climate models.

Plain Language Summary

In situ observations showed that mixed-phase clouds might exist in the form of two extremes: (a) genuinely mixed, when supercooled droplets and ice particles are uniformly distributed in a cloud volume, and (b) conditionally mixed, when ice and liquid phases are spatially separated. The objective of this study is to explore spatial scales of genuinely mixed and single-phase ice and liquid cloud segments. It was found that the spatial scales of genuine mixed-phase, ice and liquid phase clouds may vary from 100 km down to 100 m or even smaller. The obtained results are of great importance for improvements in the subgrid presentation of mixed-phase clouds in numerical weather and climate models and interpretation of remote sensing measurements.

1. Introduction

Mixed-phase clouds represent a three-phase colloidal system consisting of water vapor, ice particles, and supercooled liquid droplets. Mixed-phase clouds are ubiquitous in the troposphere, occurring at all latitudes from the polar regions to the tropics (e.g., D’Alessandro et al., 2019; Wang et al., 2013). Because of their widespread nature, mixed-phase clouds play important roles in precipitation formation, and the radiative energy balance on both regional and global scales.

One of the important characteristics of mixed-phase clouds is the degree of homogeneity of mixing ice particles and liquid droplets. There are two possible extremes of mixing. The first one is when ice particles are uniformly mixed (Figure 1a). The second one is when ice particles and liquid droplets are clustered in single-phase liquid or ice cloud regions with a complex morphology (Figure 1b). The first type of mixed-phase clouds is referred to as “genuine” mixed-phase and the second type is “conditional” mixed-phase (A. Korolev et al., 2017).

In genuine mixed-phase clouds, ice particles and liquid droplets are interacting with each other through the molecular diffusion of water vapor. In absence of dynamic forcing genuine mixed-phase clouds are colloidally unstable enabling the Wegener-Bergeron-Findeisen (WBF) process (Bergeron, 1935; Findeisen, 1938; Wegener, 1911), which results in the growth of ice particles and evaporation of droplets, and eventually leads to the glaciation of a mixed-phase cloud. Theoretical analysis suggests that glaciation time depending on temperature, ice particle concentration, and liquid water content (LWC) may vary from few minutes to tens of minutes (Korolev & Mazin, 2003). Therefore, the WBF process is an important factor limiting the endurance of genuine mixed-phase clouds.

However, in conditional mixed-phase clouds, the interaction between ice crystals and liquid droplets is hindered because of their spatial separation (Figure 1b). Therefore, under the assumption of inhibited turbulent mixing between ice and liquid cloud regions, the WBF process in such clouds will be suppressed, and conditional mixed-phase clouds would be colloidally stable. The endurance of such clouds is determined by processes other than the WBF mechanism. Since the WBF process enhances the growth rate of ice particles, then for the same
amount of ice and liquid in the same volume the rate of precipitation formation in conditional mixed-phase clouds will be slower compared to that in genuine mixed-phase clouds.

Radiative properties of genuine and conditional mixed-phase clouds are also different due to the spatial clustering of ice and liquid phase (e.g., Ruiz-Donoso et al., 2020). The liquid water, initially distributed among a large number of liquid droplets, by the end of glaciation will be depleted by fewer ice particles. Therefore, the WBF process will result in a reduction of the extinction coefficient, optical thinning, cloud coverage and lifetime of genuine mixed-phase clouds. In contrast to the above, in conditional mixed-phase, the glaciation process is hindered, and therefore, the radiative properties will be more stable compared to the genuine mixed-phase clouds. Therefore, the net radiative effect for genuine and conditional mixed-phase clouds integrated over time will be different.

Understanding the spatial phase intermittency of clouds has great importance both for fundamental cloud physics and numerical simulations of clouds. Several climate and numerical weather prediction (NWP) models use representations of cloud fractions for liquid and ice contents that are then used to determine the in-cloud hydrometeor values, from which microphysical process rates are computed (e.g., Bogenschutz et al., 2012; Walters et al., 2019). However, the degree of overlap of the in-cloud liquid and ice within grid cells, that is, the mixed-phase cloud fraction, is unconstrained. The type of mixed-phase cloudiness that is represented will impact the calculations of cloud optical properties and the WBF process rate, and thus the rate of glaciation. This also depends on the horizontal grid spacing which, depending on modeling system, varies widely. For example, NWP systems use grid spacings ranging from approximately 3 km to 15 km (e.g., McTaggart-Cowan et al., 2019), with some experiment systems on the hectometer scale (e.g., Joe et al., 2018), whereas global climate models run with grid spacings ranging from 60 to 300 km (Almazroui et al., 2020). The assumption about homogeneously mixed liquid droplets and ice particles may result in biases in precipitation production and radiation transfer.

In situ observations showed high spatial phase variability of clouds. Thus, Hallett (1999) observed well-separated regions of ice and liquid in sea-breeze-front clouds with a mixed-phase interface of only a few hundreds of meters wide. A. V. Korolev et al. (2003) showed that clusters of ice and liquid phase clouds can exist at the scale of the order of kilometers. Field et al. (2004) obtained statistics of contiguous ice, liquid, and mixed-phase cloud segments from in situ observations in frontal clouds. They pointed out that liquid intrusion into surrounding ice clouds can be as small as 100 m. A. V. Korolev et al. (2003) hypothesized that phase intermittency may exist down to a meter, that is, fine-scale clusters of liquid (ice) embedded in a glaciated (liquid) cloud.

This study is aimed to address the following questions: (a) what are the typical spatial scales of genuinely mixed-phase and single-phase ice, liquid clouds? (b) what is the spatial intermittency of single-phase ice, liquid, and genuinely mixed-phase cloud segments?

2. Methodology and the Data Set

The cloud phase composition was studied with a set of airborne instruments installed by Environment and Climate Change Canada (ECCC) in collaboration with the National Research Council (NRC) on the NRC Convair-580 research aircraft.

The phase composition of clouds was identified with the help of a set of instruments: Nevzorov probe (A. V. Korolev et al., 1998), Rosemount Icing Detector (Baumgardner & Rodi, 1989; Mazin et al., 2001), Forward Scattering Spectrometer Probe (FSSP) (Knollenberg, 1976; McFarquhar et al., 2017), Optical Array Probe 2DC (OAP-2DC) (Baumgardner et al., 2017; Knollenberg, 1976), and Optical Array Probe 2DP (OAP-2DP) (Knollenberg, 1976).

The Nevzorov probe was primarily used for the assessment of liquid water (LWC) and ice water content (IWC). The methodology of Nevzorov probe data processing and phase discrimination was described in detail in (Korolev & Strapp, 2002; A. V. Korolev et al., 1998). The Nevzorov probe liquid water sensor measurements were corrected on the residual effect of ice (Field et al., 2004; A. V. Korolev et al., 1998, 2003) and the total water sensor measurements were corrected on the ice bouncing effect (A. V. Korolev et al., 2013). The Rosemount Icing Detector was used to identify the presence of the liquid phase and exclude false liquid signals in ice clouds. The FSSP was employed to identify the presence of liquid droplets smaller than 45 μm in diameter. The OAP-2DC and 2DP were used for justification of the presence or absence of ice particles based on identification of non-circular shapes of their binary images.
In the present study the thresholds for liquid water content and ice water content (IWC) were set as LWC > 0.01 g m\(^{-3}\), IWC > 0.01 g m\(^{-3}\), respectively. The phase composition of clouds was identified based on the assessment of the ice water fraction \(\mu = \frac{\text{IWC}}{\text{LWC} + \text{IWC}}\). Thus, clouds with \(\mu > 0.9\) were considered as ice, clouds with \(\mu < 0.1\) were defined as liquid, and clouds \(0.1 \leq \mu \leq 0.9\) were determined as mixed-phase clouds. The smallest spatial scale was limited by 100 m, and it was determined by the temporal resolution of the Nevzorov probe (~1 s) and average aircraft sampling speed (~100 m s\(^{-1}\)).

The objective of the data processing was the identification of continuous ice, liquid, and mixed-phase cloud segments along the flight path of the research aircraft. Following the above definitions of clouds and cloud phase, the time series of LWC and IWC were converted into three discrete cloud phase categories of “ice,” “liquid,” and “mixed-phase.” The processing algorithm started counting the cloud length when LWC or IWC exceeded the threshold value 0.01 g m\(^{-3}\), and the cloud length counting was interrupted, when the cloud phase changed, or LWC and IWC became lower the predetermined cloud water content threshold. Isolated single-phase ice or liquid clouds were defined as continuous clouds, which were not interrupted by cloud segments with another phase. Such clouds were excluded from the subsequent analysis. The mixed-phase clouds at the smallest resolved scale (100 m) were assumed as genuinely mixed. Therefore, the ensemble of clouds used in this study consists of continuous conditional and genuine mixed-phase cloud segments.

The analyzed data set includes seven flight campaigns lead by ECCC and extended over a period of 10 years from 1994 to 2004: Beaufort Arctic Storm Experiment (BASE), Canadian Freezing Drizzle Experiment (Phase 1 and 3) (CFDE 1 and 3), the First International Satellite Cloud Climatology Project (ISCCP) Regional Experiment Arctic Cloud Experiment (FIRE-ACE), Alliance Icing Research Study (Phase 1, 1.5, 2) (AIRS 1, 1.5, 2). The description of the time frames, map of areas of operations, number of flights, and sampled cloud lengths for each campaign are available from Table S1 and Figure S1 in Supporting Information.

Most of the data was collected in stratiform clouds associated with mesoscale frontal systems. The altitude and temperature of the sampled clouds varied in the ranges 0.5 km < \(H\) < 7.3 km and −50 < \(T\) < +10°C, respectively. However, in the frame of this study the measurements were limited by the temperature range −35 < \(T\) < 0°C due to the low statistics of mixed and liquid clouds in the range −40 < \(T\) < −35°C, and the absence of liquid phase at \(T\) < −40°C due to homogeneous freezing.

The entire flight operations include 117 research flights. The endurance of each flight varied from 3 to 5 hr. The total length of sampled in-cloud space extended over 55,381 km. After elimination of isolated single-phase clouds and applying temperature limitations, the total length of the clouds, which were included in the statistics, was reduced to 32,488 km.

### 3. Results

Figure 2 shows frequency distributions of ice, liquid, and mixed-phase contiguous cloud segments in three temperature subranges −10 < \(T\) < 0°C, −20 < \(T\) < −10°C, and −30 < \(T\) < −20°C. The distributions were...
normalized on the total length of clouds in each phase category to facilitate their intercomparisons with each other. As it is seen from Figure 2 the ice, liquid, and mixed-phase cloud regions are represented by a cascade of scales ranging from tens and hundreds of kilometers down to the minimum scale \( L_{\text{min}} \approx 100 \) m. The minimum scale \( L_{\text{min}} \) is limited by the instrumental resolution.

The maximum scale of ice clouds \( (L_{\text{max}})^{(i)} \) varies from tens kilometers at cold temperatures \((-30 < T < -20{\degree}\text{C})\) (Figure 2c) to 100–300 km at \( T > -20{\degree}\text{C} \) (Figures 2a and 2b). The maximum scales of liquid \( (L_{\text{max}})^{(w)} \) and mixed-phase \( (L_{\text{max}})^{(m)} \) cloud segments are systematically lower than \( (L_{\text{max}})^{(i)} \) for ice clouds.

As it is seen from Figure 2, \( (L_{\text{max}})^{(w)} \) and \( (L_{\text{max}})^{(m)} \) decrease from 40 to 60 km at \(-10 < T < 0{\degree}\text{C}\) to 10–15 km at \(-30 < T < -20{\degree}\text{C}\).

The spatial scale distributions of continuous ice, liquid, and mixed-phase segments are well described by the power law \( (L) = aL^{b} \). The weighted least squares fit coefficients \( a \) and \( b \) are presented in Table 1. The Pearson correlation coefficient for each parameterization in Table 1 are higher 0.99.

As it is seen from Table 1 the slope of lengths distributions for mixed-phase and liquid (Figure 2) remains nearly constant in all temperature intervals, and it changes within 3%–4%. However, the slope of lengths distributions for ice has a clear tendency to increase toward low temperatures.

Figure 3 shows temperature dependences of average continuous lengths of ice, liquid, and mixed phase cloud segments. As it is seen from Figure 3 the average length of mixed-phase segments has a weak dependence on temperature and it varies between 300 and 500 m. The mean length of liquid cloud segments is systematically larger than that of mixed-phase segments, and it gradually decreases with the decrease of temperature from approximately 900 m down to 500 m. However, average lengths of continuous ice cloud segments turned out to be the largest compared to liquid and mixed-phase segments. The average length of ice cloud segments increases from 1 km at \( 0{\degree}\text{C} \) to 5 km at \(-25 < T < -20{\degree}\text{C}\) and then decrease to approximately 1 km at \(-35{\degree}\text{C}\).

4. Discussion

The slopes of obtained frequency distributions or some of their parts in Figure 2 are close to \(-5/3\). In the atmosphere, the \(-5/3\) power law is underlying numerous processes, including turbulence, storm intensity, ocean wave growth, etc. The proximity of distributions of cloud spatial scales to the universal \(-5/3\) power law is indicative of the linkage between formation of mixed-phase clouds and turbulence.

Turbulence is one of the major mechanisms mixing environments with different phases. The result of the turbulent mixing of clouds with a different phase composition is a mixed-phase cloud. The characteristic time scale \( (\tau) \) of mixing of a cloud volume with a spatial scale \( L \) can be assessed as (Landau & Lifshitz, 1987)

\[
\tau = \left( \frac{L^2}{\epsilon} \right)^{1/3}
\]

where \( \epsilon \) is the turbulence energy dissipation rate.

However, cloud particles are an active admixture and during mixing they are interacting with each other through water vapor. Such interaction results in activation of the WBF mechanism, if the vertical velocity of the cloud \( u_z \)
Substituting Equation 1 and 2 in Equation 3 yields a threshold spatial scale,

$$L_{th} = (k^3 \varepsilon)^{1/3} \frac{W}{N_i}$$  \hspace{1cm}  (4)

such that mixed-phase environment can exist at scales $L < L_{th}$ at time scales satisfying Equation 3.

Substituting $\varepsilon = 10^{-3} - 10^{-4}$ m$^2$ s$^{-1}$, $W = 0.01 - 1.0$ g m$^{-3}$, $N_i = 10^3 - 10^4$ L$^{-1}$, $-35 < T < -5$°C yields the range of spatial phase scales $L_{th} \sim 10^3 - 10^4$ m. This obtained assessment of $L_{th}$ remarkably consistent with the range of spatial scales covered by mixed-phase clouds in Figure 2. However, the obtained assessment, also suggests that $L_{th}$ may extend toward the small-scale end. As it is seen from Equation 4, the spatial scale of mixed-phase clouds depends on LWC. Therefore, it is anticipated that the frequency function $F(L)$ should be related on a frequency of occurrence of LWC ($F(W)$). Figure 4 shows a probability density functions $F(W)$ at different temperatures. As it is seen from Figure 4 $F(W)$ increases with the decrease of LWC. Therefore, following Equation 4 the increase of occurrence of small-scale mixed-phase clouds may be explained by the increase of occurrence of LWC with the decrease of LWC. Even though $F(W)$ has a quasi-exponential dependence, whereas $F(L)$ is described by a power-law, the behavior of LWC generally explains the behavior of $F(L)$.

It should be noted that the above simplified consideration is relevant for isotropic turbulence in the inertial subrange, which in free atmosphere is limited by approximately 1 km (Honnert et al., 2020; Wyngaard, 2010). The spatial scales of the cloud phase composition at larger scales are controlled by convection and mesoscale dynamics (Wood & Field, 2011). The above consideration also did not account for cloud processes such as riming, ice sedimentation, and entrainment of out-of-cloud air. The later was in detail discussed in Pinsky et al. (2018). However, the above approach allows to generally predict the behavior and the range of spatial scales of genuinely mixed-phase clouds.

In the present study the smallest spatial scale of genuine mixed-phase cloud segments was limited by the instrumental resolution (100 m). In fact, these clouds can be conditionally mixed at spatial scales $L < L_{min}$. It should be

Table 1: Coefficients for the Power Law Fitting of the Normalized Frequency Distributions of Spatial Scales of Ice, Liquid, and Mixed-Phase Cloud Regions $F(L) = aL^b$

| Cloud type | $-30 < T < -20^\circ$C | $-20 < T < -10^\circ$C | $-10 < T < 0^\circ$C |
|------------|-------------------------|-------------------------|-----------------------|
| Ice        | $a = 10.5 \times 10^{-3}$, $b = -1.44$ | $a = 18 \times 10^{-3}$, $b = -1.66$ | $a = 17.4 \times 10^{-3}$, $b = -1.86$ |
| Mixed      | $a = 7.85 \times 10^{-3}$, $b = -2.03$ | $a = 8.06 \times 10^{-3}$, $b = -2.17$ | $a = 8.75 \times 10^{-3}$, $b = -2.19$ |
| Liquid     | $a = 24.4 \times 10^{-3}$, $b = -1.97$ | $a = 11.5 \times 10^{-3}$, $b = -2.03$ | $a = 13.8 \times 10^{-3}$, $b = -1.94$ |

Figure 3: Average lengths of ice, liquid, and mixed-phase cloud segments versus temperature.
is the average distance between particles, for ice particle and liquid droplet clusters of the genuine conditionally ∼ 0.01–10 L^{-1} for ice may go down to a centimeter scale. However, is the average ∼ 10–1,000 cm^{-3} for typical ice particle concentrations \( n_{\text{tot}} \sim 0.01–10 \) L^{-1} the minimum scale of ice is expected to be 0.1 m < \( L_{\text{min}} < 1 \) m. For the cases of ice multiplication (i.e., \( N_i \sim 10^3 \) L^{-1}) \( L_{\text{min}} \) for ice may go down to a centimeter scale. However, for cloud droplet at concentration \( N_d \sim 10–1,000 \) cm^{-3} the minimum cluster scale is estimated to vary in the range 2 mm < \( L_{\text{min}} < 10 \) mm.

As seen, the obtained estimates of \( L_{\text{min}} \) for ice particle and liquid droplet clusters are much smaller than that in this study. The question of \( L_{\text{min}} \) is of great importance, and its future addressing requires advanced in situ observations.

5. Conclusions

This study presents observed spatial phase intermittency in mixed phase clouds based on a large data set collected in midlatitude and Arctic clouds. It was shown that in the temperature range −35 < \( T < 0 \)°C clouds containing ice and liquid may form regions with genuinely or conditionally mixed-phase environments. The lengths of continuous segments of ice, liquid, and mixed-phase cloud present a cascade of scales varying from 10 km down to a minimum scale \( L_{\text{min}} \sim 100 \) m determined by the spatial resolution of the measurements. It is hypothesized that \( L_{\text{min}} \) of the genuine conditionally mixed-phase cloud regions may go down to scale of 1–10 m or even smaller.

The results obtained yield insight on the spatial morphology of mixed-phase clouds. These results also suggest the potential of using in situ observations to validate and calibrate atmospheric models in terms their ability to represent subgrid-scale phase heterogeneity. They also suggest that even in cloud-resolving and high-resolution mesoscale models (with horizontal grid spacings on the order of a few km or less) whose microphysics schemes typically do not consider cloud fractions, subgrid-scale heterogeneity of mixed-phase clouds is likely not well represented, biased toward genuinely mixed, and is a potential source of model error. Improving the type of mixing for mixed-phase clouds in models of all types will improve the calculation of cloud optical properties, thereby improving radiative transfer calculations, and modify the WBF process, which ultimately determines the glaciation time and impacts the production of precipitation and cloud brightness.

Data Availability Statement

Data supporting this research are available from https://doi.org/10.5281/zenodo.6558498.

Acknowledgments

The BASE project was funded by ECCC. The CFDE project was supported by Canadian Search and Rescue, Transport Canada (TC). The FIRE-ACE project was supported by NASA. The AIRS project was supported by the US Federal Aviation Administration (FAA) and TC. Special thanks for the NRC Convair-580 pilots for their great cooperation in data collection, and to the ECCC and NRC teams of technicians and engineers for installation, maintenance, and operations of the airborne instrumentation.

References

Almarouzi, M., Saeed, S., Saeed, F., Islam, M. N., & Ismail, M. (2020). Projections of precipitation and temperature over the South Asian countries in CMIP5. Earth Systems and Environment, 4(2), 297–320. https://doi.org/10.1007/s41748-020-00157-7
Baumgardner, D., Abel, S. J., Axisa, D., Cotton, R., Crosier, J., Field, P., et al. (2017). Cloud ice properties: In situ measurement challenges. Meteorological Monographs, 58, 9.1–9.23. https://doi.org/10.1175/AMSMONOGRAPH-D-16-0011.1
Baumgardner, D., & Rodi, A. (1989). Laboratory and wind tunnel evaluation of the Rosemount icing detector. Journal of Atmospheric and Oceanic Technology, 6, 971–979. https://doi.org/10.1175/1520-0426(1989)006<0971:LMIVLT>2.0.CO;2
Bergeron, T. (1935). On the physics of clouds and precipitation. In Proces Verbaux de l’Association de Météorologie (pp. 156–178). International Union of Geodesy and Geophysics.
Bogenschutz, P. A., Gettelman, A., Morrison, H., Larson, V. E., Schanen, D. P., Meyer, N. R., & Craig, C. (2012). Unified parameterization of the planetary boundary layer and shallow convection with a higher-order turbulence closure in the community atmosphere model: Single-column experiments. Geoscientific Model Development, 5(6), 1407–1423. https://doi.org/10.5194/gmd-5-1407-2012
D’Alessandro, J. J., Diao, M., Wu, C., Liu, X., Jensen, J. B., & Stephens, B. B. (2019). Cloud phase and relative humidity distributions over the Southern Ocean in austral summer based on in situ observations and CAM5 simulations. Journal of Climate, 32(10), 2781–2805. https://doi.org/10.1175/JCLI-D-18-0232.1

Figure 4. Normalized frequency of occurrences of liquid water content in liquid and mixed-phase clouds averaged over all sampled clouds. Isolated single-phase clouds were not included in these statistics.
Field, P. R., Hogan, R. J., Brown, P. R. A., Illingworth, A. J., Choularton, T. W., Kaye, P. H., et al. (2004). Simultaneous radar and aircraft observations of mixed-phase cloud at the 100 m scale. Quarterly Journal of the Royal Meteorological Society, 130(600), 1877–1904. https://doi.org/10.1256/qj.03.102

Findeisen, W. (1938). Kolloid-meteorologische Vorgänge bei Neiderschlags-bildung. Meteorologische Zeitschrift, 55, 121–133. https://doi.org/10.1127/metz/2015/00675

Hallett, J. (1999). Charge generation with and without secondary ice production. In Proceedings of 11th international conference on atmospheric electricity (pp. 355–358). Huntsville, NASA/CP-1999-209261, June.

Honmet, R., Efstatiou, G., Beare, R., Ito, J., Lock, A., Neggers, R., et al. (2020). The atmospheric boundary layer and the “gray zone” of turbulence: A critical review. Journal of Geophysical Research: Atmospheres, 125(13), e2019JD030317. https://doi.org/10.1029/2019JD030317

Joe, P., Belair, S., Bernier, N. B., Bouchet, V., Brook, J. R., Brunet, D., et al. (2018). The environment Canada pan and parapan American science showcase project. Bulletin of the American Meteorological Society, 99(5), 921–953. https://doi.org/10.1175/BAMS-D-16-0162.1

Knollenberg, R. G. (1976). Three new instruments for cloud physics measurements: The 2D-spectrometer probe, the forward scattering spectrometer probe and the active scattering spectrometer probe. Preprints. In International conference on cloud physics (pp. 554–561). American Meteorological Society.

Korolev, A. (2007). Limitations of the Wegener-Bergeron-Findeisen mechanism in the evolution of mixed-phase clouds. Journal of the Atmospheric Sciences, 64, 3372–3375. https://doi.org/10.1175/JAS4035.1

Korolev, A., McFarquhar, G., Field, P. R., Franklin, C., Lawson, P., Wang, Z., et al. (2017). Mixed-phase clouds: Progress and challenges. Meteorological Monographs, 58, 1–5.50. https://doi.org/10.1175/AMSMONOGRAPHS-D-17-0001.1

Korolev, A. V., Isaac, G. A., Cober, S., Strapp, J. W., & Hallett, J. (2003). Microphysical characterization of mixed-phase clouds. Quarterly Journal of the Royal Meteorological Society, 129(587), 39–66. https://doi.org/10.1256/qj.01.204

Korolev, A. V., & Mazin, I. P. (2003). Supersaturation of water vapor in clouds. Microphysical processes and interface control document. Geosciences Monograph, 600, 1877–1904.

Korolev, A. V., Strapp, J. W., Isaac, G. A., & Nevzorov, A. N. (1998). The Nevzorov airborne hot-wire LWC–TWC probe: Principle of operation and performance characteristics. Journal of Atmospheric and Oceanic Technology, 15(6), 1495–1510. https://doi.org/10.1175/1520-0426(1998)015<1495:TNAHWL>2.0.CO;2

Laudau, L. D., & Lifshitz, E. M. (1987). Fluid mechanics (Vol. 6, p. 539). Pergamon Press.

Mazzin, I. P., Korolev, A. V., Heymsfield, A., Isaac, G. A., & Cober, S. G. (2001). Thermodynamics of icing cylinder for measurements of liquid water content in supercooled clouds. Journal of Atmospheric and Oceanic Technology, 18(4), 543–558. https://doi.org/10.1175/1520-0426(2001)018<0543:TOICFM>2.0.CO;2

McFarquhar, G. M., Baumgardner, D., Baisaner, A., Abel, S. J., Crosier, J., French, J., et al. (2017). Processing of ice cloud in situ data collected by bulk water, scattering, and imaging probes: Fundamentals, uncertainties and efforts towards consistency. Meteorological Monographs, 58, 1–11. https://doi.org/10.1175/AMSMONOGRAPHS-D-16-0007.1

McTaggart-Cowan, R., Vaillancourt, P. A., Zadra, A., Chamberland, S., Charron, M., Corvec, S., et al. (2019). Modernization of atmospheric physics parameterization in Canadian NWP. Journal of Advances in Modeling Earth Systems, 11(11), 3593–3635. https://doi.org/10.1029/2019ms001781

Pinsky, M., Khain, A., & Korolev, A. (2018). Theoretical analysis of liquid–ice interaction in the unsaturated environment with application to the problem of homogeneous mixing. Journal of the Atmospheric Sciences, 75(4), 1045–1062. https://doi.org/10.1175/JAS-D-17-0228.1

Ruíz-Donoso, E., Ehrlich, A., Schäfer, M., Jakell, E., Schermann, V., Crowell, S., et al. (2020). Small-scale structure of thermodynamic phase in Arctic mixed-phase clouds observed by airborne remote sensing during a cold air outbreak and a warm air advection event. Atmospheric Chemistry and Physics, 20(9), 5487–5511. https://doi.org/10.5194/acp-20-5487-2020

Walters, D., Baran, A. J., Boutle, I., Brooks, M., Earnshaw, P., Edwards, J., et al. (2019). The met office unified model global atmosphere 7.0/7.1 and JULIE global land 7.0 configurations. Geoscientific Model Development, 12(5), 1909–1963. https://doi.org/10.5194/gmd-12-1909-2019

Wang, Z., Vane, D., Stephens, G., & Reinke, D. (2013). CloudSat Project: Level 2 combined radar and lidar cloud scenario classification product process description and interface control document (p. 61). California Institute of Technology JPL. Retrieved from https://www.cloudsat.cira.colostate.edu/cloudsat-static/infold/2b-cldclass-lidar/2B-CLDCLASS-LIDAR_PDI(IDC_P_R04.20120522.pdf

Wegener, A. (1911). Thermodynamik der atmosphäre (p. 331).

Wood, R., & Field, P. R. (2011). The distribution of cloud horizontal sizes. Journal of Climate, 24(18), 4800–4816. https://doi.org/10.1175/2011JCLI4056.1

Wyngaard, J. C. (2010). Turbulence in the atmosphere (p. 406). Cambridge University Press.