Optimizing the Frequency of Quantum Error Correction using the [[7,1,3]] Steane Code

Ali Abu-Nada\textsuperscript{1}, Ben Fortescue\textsuperscript{1}, Mark Byrd\textsuperscript{1,2}

\textsuperscript{1}Physics Department, Southern Illinois University, Carbondale, IL 62901, USA and
\textsuperscript{2}Computer Science Department, Southern Illinois University, Carbondale, IL 62901, USA
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A common assumption in analyses of error thresholds and quantum computing in general is that one applies fault-tolerant quantum error correction (FTQEC) after every gate. This, however, is known not to always be optimal if the FTQEC procedure itself can introduce errors. We investigate the effect of varying the number of logical gates between FTQEC operations, and in particular the case where failure of a postselection condition in FTQEC may cause FTQEC to be skipped with high probability.

By using a simplified model of errors induced in FTQEC, we derive an expression for the logical error rate as a function of error-correction frequency, and show that in this model the optimal frequency is relatively insensitive to postselection failure probability for a large range of such probabilities. We compare the model to data derived from Monte Carlo simulation for the [[7, 1, 3]] Steane code.

I. INTRODUCTION

Quantum error correcting codes (QECC) are often required to correct errors due to imperfectly implemented quantum operations during quantum information processing and also to simply store quantum information\cite{Kliesch}. For a quantum error correcting code to be effective, the probability for an error to occur on any given gate needs to be below the threshold value for the code\cite{Knill}. However, even if the error rate is below the threshold, the resources required for quantum error correction can greatly exceed those resources required for the quantum information processing alone.

QECCs encode logical qubits into multiple physical qubits; additional quantum error correction (QEC) operations are regularly performed to diagnose and repair errors that may have arisen. (In this paper we use QEC to denote the operations required to obtain an error syndrome from data and correct any detected errors, although in the examples we consider the latter does not require any additional physical operations).

A common model for such an error-corrected computation is to apply a QEC immediately after each logical gate in order to correct the physical errors introduced by these gates before they accrue into more serious logical errors. However, QEC operations will, in general, also introduce errors with some non-zero probability since they also use imperfect physical gates. These are often the same gates used to perform logical operations. In addition to the threshold needing to be satisfied, for QEC operations to successfully suppress errors enough to allow large-scale computation, the QEC as well as logical gates must be implemented in a fault-tolerant way\cite{Knill}, i.e., in a way such that a single faulty quantum gate cannot lead to multiple errors on the data. Fault-tolerant constructions for QEC operations are often dependent on post-selection; in particular, they commonly use ancillary states to diagnose errors, which effectively carry away entropy from the data. Since they interact with the data, these states must be prepared so that they do not spread multiple errors to the data. Often, however, non-fault-tolerant circuits must be used for the initial ancilla preparation, and fault-tolerance is instead enforced by post-selection (ancilla verification) in which only those ancillas which satisfy some measurement outcome after being created are subsequently used to interact with the data\cite{Knill}.

This model of fault-tolerant QEC, in which ancillas are created until one passes post-selection, can be difficult to implement. The data can accumulate additional errors either waiting for a “good” ancilla to be created (if created sequentially) or to be moved into place (if created in parallel). And the nondeterministic delays involved make synchronization of the data with other data blocks difficult. An alternative, which preserves synchronization, is to simply to allow for a fixed number of postselection attempts, and to carry on with the computation (skipping the QEC) if these are unsuccessful. The obvious disadvantage is that skipped QEC operations allow errors to accumulate from sequential logical gates, but if the skipping probability and gate error probabilities are sufficiently low this may still be the optimal solution. Another alternative is go through with the error correction without post-selection and instead measure to detect any errors on the ancilla and correct for those as well as any error already present. This is called ancilla decoding and was originally shown to be advantageous in the regime of slow, noisy measurements\cite{Weinstein}, but has recently been considered in other contexts\cite{Cohen}. A closely related situation is where QEC operations are intentionally not applied after every gate, but only after a certain number of gates. This can reduce the overall logical error if the errors introduced by the QEC operation are larger than those introduced by the gates. Recently, Weinstein has provided a specific relation between the fidelity and physical error rates for performing QECs after different numbers of gates in the [[7, 1, 3]] Steane code\cite{Weinstein}. Here, it is shown that the overall error rate is
not significantly increased by performing QEC after two gates rather than after every gate, and that the inherently noisy correction process will sometimes introduce more error into the system when applied too frequently [9, 10].

In this paper, we consider the optimal frequency to apply postselection-dependent QEC operations, in the case where postselection failure results in a skipped QEC. We consider a model where a logical data qubit undergoes operations from $N$ logical gates, with $m$ gates in between each QEC, using the well-known [[7,1,3]] Steane code $[3]$ and the Steane ancilla technique for the latter. We determine the logical error rate $P_L$ for the data after undergoing these $(N/m)$ “blocks” of gates and QEC operations, as a function of $m$ and the underlying physical gate error $\epsilon$, and then minimize as a function of $m$.

In section II we explain the model and the analytical derivation of the formula given the model. In section III we compare the predicted error rate of the model to that obtained via Monte Carlo simulation.

II. DERIVATION OF THE MATHEMATICAL FORMULA FOR $P_L$

A. Logical gate and QEC model

In general, the dependence of the logical error rate on the physical error rate of the underlying gates depends on the circuits used to implement gates and QEC. These circuits can be complex, and highly dependent on the chosen code. While our analysis is based on the Steane code, we wished to use a model that could be readily adapted to other codes. We thus produced a semi-abstracted model based on the Steane code with Steane ancillas.

Errors may be introduced into the logical qubit in two ways, from the logical gates and the “noisy” QEC itself (in this analysis we do not treat errors from movement or hold operations as a separate category, they may be included into the above categories if desired). We model a noisy physical gate as performing the desired operation followed by, with probability $\epsilon_g$, an error operation. We treat logical gates as transverse, that is, simply consisting of a single physical gate applied to each qubit.

Our model of the QEC operation is more approximate. We divide errors induced by QEC into four separate parts (as shown in Figure 1) with the following probabilities:

- “Correction errors”, with probability $\epsilon_o$ per qubit, are due to physical errors in those gates directly applied to the data as part of the QEC process. In the Steane code and many others, this is limited to those two-qubit gates used to interact the data with the ancillas (since the corrections themselves can be done using the “Pauli frame” [11], without the use of physical gates). We defined correction errors as being those which affect the data only, not the QEC ancilla. Hence two-qubit gate failures which lead to errors on both outputs (both data and ancilla qubits) are excluded.

- The remaining errors are those which affect the data through causing an incorrect syndrome measurement, and thus an incorrect correction operation. We first define “syndrome errors”, with probability $\epsilon_s$ per qubit, as those errors where an error on the QEC ancilla (or its measurement) only (with no errors directly applied to the data) cause a data qubit to be wrongly “corrected”, in addition to any errors already present.

- “Omission errors”, with probability $\epsilon_o$ per qubit, represent the special class of erroneous syndromes which coincidently combine with existing data errors (where present) to (wrongly) return a syndrome indicating no errors. Thus if the data is initially without error, an omission error in the QEC will lead to an error on the data, but if a single qubit error is present on the data prior to the QEC, an omission error will simply lead to this error remaining uncorrected.

- Finally, “double errors”, with probability $\epsilon_d$ per qubit, are those errors where a single gate failure in the CNOT joining the data to the ancilla lead to X errors on both the source and the target (thus we only need consider this class of errors when modeling source and target errors on two-qubit gates as correlated). This behaves differently to a standard syndrome error since if a double error occurs during syndrome measurement, is the only error occurring during that syndrome measurement, and the data has no prior errors, no error will be added to the data from the overall QEC. This is because an error on both source and target qubits in the CNOT is equivalent to an error on the data which correctly propagates to the syndrome, and will therefore be successfully corrected. Conversely, since a double error does affect the data directly, if a double error on occurs on one qubit and a syndrome (or double) error on another qubit as part of the same syndrome measurement, this can lead to two errors on the data (while two standard syndrome errors could only lead to one error on the data).

We additionally consider the probability of an “ancilla error”, with probability $\epsilon_a$ (not per qubit). This is the case where a QEC operation (that is, any elements of that operation that could either correct errors or produce errors on the data) is not performed (for example, due to failure in the postselection process for the necessary ancilla). Such events do not produce data errors themselves, but result in existing errors not being corrected when they should be.

Note that all of the above are functions of the physical gate error $\epsilon_g$, but the exact relationship depends on the circuits and QECC used, thus we treat them as separate variables in our model.
B. Logical error rate

The sequence of \((N/m)\) blocks may, in a distance-3 code, produce a logical error if one or more of these blocks create two or more physical errors on the logical qubit. Therefore, to estimate the logical error rate \(P_L\), we must enumerate the ways in which these blocks might create two or more physical errors on the logical qubit. This logical error might be created either within a single block, or across multiple blocks (if a QEC is either skipped due an ancilla error or fails due to a syndrome error.) Thus a logical error will only occur with probability second-order or higher in the various error probabilities.

We are particularly interested, however, in the regime where the ancilla errors are significantly larger than the other errors. This can easily be the case if ancilla creation circuits are complex. While the postselection procedure means such large errors do not translate directly to logical errors, they can result in multiple QEC operations being skipped, with a consequent increase in the logical error probability if other errors are present.

Considering now the structure of a block, this consists of \(m\) gates, followed by a QEC operation consisting of a successful correction of any errors on the logical qubit, unless an error occurs in the QEC. Finally any correction errors are applied. Thus, in the absence of QEC verification failures, the leading-order contributions to the logical error rate (where two errors occur resulting in a logical error) are very limited: two errors can occur within a block or across two adjacent blocks, as shown in detail in Tables II to IV.

A verification failure permits additional error combinations: errors on two separate blocks between which a successful correction would ordinarily occur. Note that additional verification failures permit, to second order, the same types of logical errors. That is, the only second-order errors which \(f\) skipped QECs additionally permit (besides those which could occur regardless) are when the skipped QECs all occur sequentially, and the two errors in question are on the block containing the first skipped QEC and the block following the final skipped QEC. For a sequence of \(B \equiv N/m\) blocks, there are \(B - f\) ways to have \(f\) sequential skipped QECs. Thus the additional logical error due to \(f\) skipped QECs is weighted by a factor

\[
\gamma \equiv \sum_{f=1}^{B-1} \epsilon_a^f (B - f) \\
= B\epsilon_a(1 - \epsilon_a) - \epsilon_a + \epsilon_a^{B+1} \\
= \frac{B\epsilon_a(1 - \epsilon_a) - \epsilon_a + \epsilon_a^{B+1}}{(1 - \epsilon_a)^2}
\]

as used in Table III. Due to the definition of correction errors, in some cases the combined errors allowable due to a single skipped QEC may span 3 rather than 2 blocks, in which case the multiplying factor is

\[
\gamma_3 \equiv \sum_{f=1}^{B-1} \epsilon_a^f (B - f - 1) \\
= (\gamma - (B - 1)\epsilon_a)/\epsilon_a \\
= \frac{\epsilon_a}{(1 - \epsilon_a)^2} (B(1 - \epsilon_a) + 2\epsilon_a + \epsilon_a^{B-1})
\]

as used in Table IV. In Tables II to IV we show the possible ways for a logical error to occur to second order in the various errors. The final column represents the overall contribution to the logical error rate from all error combinations of that type. Note that rows represent a general class of errors (e.g. the first row in Table II represents the contribution due to two gate errors within a single block) and the rows involving factors of \(\gamma\) and \(\gamma_3\) represent classes of errors where \(f\) intermediate QEC steps are skipped leading to an overall factor of \(\gamma\) or \(\gamma_3\).

Summing the terms, the general second-order formula for \(P_L\) is:
\[ P_L = 42\left[B\left(me_g\left(\frac{me_g}{2} + (1 - \epsilon_a)(\epsilon_s + \epsilon_d)\right) + (1 - \epsilon_a)\left(c_c \left(\epsilon_s + \epsilon_a + \frac{\epsilon_a}{2}\right) + \epsilon_d \left(\epsilon_s + \frac{\epsilon_d}{2}\right)\right)\right)
+ (B - 1 + \gamma_3)(1 - \epsilon_a)(\epsilon_c + \epsilon_s + \epsilon_a)(m\epsilon_g + (1 - \epsilon_a)(\epsilon_s + \epsilon_d)) + \gamma m\epsilon_g (m\epsilon_g + (1 - \epsilon_a)(\epsilon_s + \epsilon_d))\right] \] (6)

C. Minimizing \( P_L(m) \)

\( P_L \) is a discrete function of \( m \). In the limit of many blocks \( (B \to \infty) \) we can neglect express \( P_L \) as a function

\[ P_{\epsilon_p}(m) \simeq dm^{-1} + c_0 + c_1m \] (7)

where

\[ d = 42B(1 - \epsilon_a)\left[e_c \left(\epsilon_s + \epsilon_a + \frac{\epsilon_a}{2}\right) + \epsilon_d \left(\epsilon_s + \frac{\epsilon_d}{2}\right)\right] + (\epsilon_s + \epsilon_d)(\epsilon_c + \epsilon_s + \epsilon_a) \] (8)

\[ c_0 = 42Bc_g(\epsilon_c + 2\epsilon_s + \epsilon_a + \epsilon_d) \] (9)

\[ c_1 = 42Bc_g^2\left(\frac{1}{1 - \epsilon_a} - \frac{1}{2}\right) \] (10)

\( \epsilon_{\text{min}} \) satisfies

\[ P_L(m_{\text{min}}) < P_L(m_{\text{min}} - 1) \]

\[ P_L(m_{\text{min}}) < P_L(m_{\text{min}} + 1), \]

\[ \frac{d}{c_1} = \frac{2(1 - \epsilon_a)^2[ec_c(\epsilon_s + \epsilon_a + \frac{\epsilon_a}{2}) + \epsilon_d(\epsilon_s + \frac{\epsilon_d}{2}) + (\epsilon_s + \epsilon_d)(\epsilon_c + \epsilon_s + \epsilon_a)]}{\epsilon_g^2(1 + \epsilon_a)} \] (12)

Thus from [7] we have

\[ m_{\text{min}}(m_{\text{min}} + 1)c_1 - d > 0 \]

\[ m_{\text{min}}(m_{\text{min}} - 1)c_1 - d > 0 \]

and hence (since \( m_{\text{min}} \) is positive)

\[ m_{\text{min}} > \sqrt{\frac{1}{4} + \frac{d}{c_1} - \frac{1}{2}} \]

\[ m_{\text{min}} < \sqrt{\frac{1}{4} + \frac{d}{c_1} + \frac{1}{2}} \] (11)

Thus the dependence of \( P_L \) on \( m \) is determined by the variable

\[ \frac{d}{c_1} = \frac{2(1 - \epsilon_a)^2[ec_c(\epsilon_s + \epsilon_a + \frac{\epsilon_a}{2}) + \epsilon_d(\epsilon_s + \frac{\epsilon_d}{2}) + (\epsilon_s + \epsilon_d)(\epsilon_c + \epsilon_s + \epsilon_a)]}{\epsilon_g^2(1 + \epsilon_a)} \] (12)

III. MONTE CARLO SIMULATION OF \( P_L \)

As discussed above, our analytical formula for the logical error simplifies the description of QEC errors (in general a function of complex ancilla circuits) to the variables \( \epsilon_{s,c,d} \), which we assume to take the same set of values for every qubit. In order to check the accuracy of this approximation, we performed Monte Carlo simulations of the complete QEC for the [[7,1,3]] Steane code with the Steane ancilla technique, using QASM-P, simulation software based on QASM [12], in order to compare the logical error rates obtained with those predicted.

Initially, all gates were simulated using the stochastic error model for depolarizing noise [8]. In this case, we considered bit-flip (X) errors only on the data qubits (phase-flip (Z) errors may be dealt with independently in the [[7,1,3]] code), and in our chosen noise model, will occur at equal rates). We used \( N = 1000 \) with varying block sizes \( m \in \{1, 2, 4, 5, 8, 10, 20, 25, 100\} \), thus \( B \) varied between 1000 and 10.

Ancilla verification is performed as usual for X correction using the Steane technique in the [[7,1,3]] code: an ancilla interacts with the data as the control for a transversal CNOT gate, and needs to be prepared and verified so that a single gate failure will not cause the ancilla to have multiple errors which would be transferred to the data. Usually, then, ancilla verification failure rates are, like the other errors, a function of the gate error rates. However, in order to vary \( \epsilon_a \) independently of other errors to verify the formula, our simulation artificially determines beforehand whether a verification failure error will occur. If so, the QEC is skipped. If not, the preparation and verification is repeated until passed. Thus verified ancillas have the correct error statistics for a given gate error, but failure occurs with the chosen probability \( \epsilon_a \).

To determine the logical error probability \( P_L \), the data is prepared, without error, in a logical |0⟩ state. Then, as
described above, a series of blocks of $m$ transversal logical gates (since we only wish to simulate errors, these are simply wait operations which in the absence of errors do not change the qubits’ state), followed by one QEC operation per block, are applied, for a total of $N$ logical gates and $B = N/m$ blocks and (attempted) QECs. Finally, the data is checked for logical $X$ errors. Each simulation (for a given choice of variable values $\epsilon_g$ and $\epsilon_d$ has $10^6$ runs). From the Monte Carlo simulation we therefore obtain $P_L$ as a function of the underlying physical error rates.

A. Numerical estimation of $\epsilon_s$ and $\epsilon_o$

By our definition, the only source of correction errors is the CNOT gate interacting the data with the ancilla. Such errors occur only when the gate failure leads to an $X$ error on the CNOT source (the data) but not the CNOT target. Similarly double errors occur when a CNOT failure leads to $X$ errors on both outputs. In our depolarizing error model single qubit gates undergo $X$, $Y$ or $Z$ errors with equal probability $\epsilon/3$, and two-qubit gates undergo the 15 possible two-qubit errors $(X \otimes I, Y \otimes I \ldots Z \otimes Z)$ with equal probability $\epsilon/15$. Since our analysis only considers bit errors (introduced by either $X$ or $Y$ operators), we have a single-qubit gate bit error probability of $\epsilon_g = 2\epsilon/3$. Thus the probability of a CNOT source-only error, which comes from the operations $X \otimes I, X \otimes Z, Y \otimes I, Y \otimes Z$ is $\epsilon_c = 4\epsilon/15 = 2\epsilon_g/5$. Likewise the probability of of a double error comes from $X \otimes X, X \otimes Y, Y \otimes X, Y \otimes Y$ and hence $\epsilon_d = 2\epsilon_g/5$

$\epsilon_s$ and $\epsilon_o$ were determined directly from simulation. By our definition, a syndrome error or double error in a QEC will, for an input logical qubit containing one error, add a second error, leading to an overall logical error, and are the only first-order QEC errors which do this. Thus to estimate error rate $\epsilon_s + \epsilon_o$, the data was first prepared in a logical eigenstate, with an error on one of the seven qubits. The QEC procedure for the [[7,1,3]] Steane code with the Steane ancilla technique was then performed using the stochastic error model for depolarizing noise. Finally the logical qubit was checked for logical errors to determine the error rate.

Similarly, if the input data has a single logical error entering and leaving the QEC, this will be due to either a correction or omission error. Hence to estimate the QEC physical error rate $\epsilon_c + \epsilon_o$, we prepare the input logical qubit with an error on one of the seven corresponding physical qubits, then we perform the same QEC simulation procedure as before, but determine the rate based on events when the output has a single error (rather than two). In both cases we varied the input error over all 7 qubits and took the mean resultant $P_L$.

We performed the simulation with a variety of numerical values for $\epsilon_g$ ($10^{-5} \sim 10^{-3}$). For each different value of $\epsilon_g$, we determine the numerical values of $\epsilon_s$ and $\epsilon_c + \epsilon_o$. The relationship is fitted to a linear equation to determine the coefficients for $\epsilon_s$ vs. $\epsilon_g$ and $\epsilon_c + \epsilon_o$ vs. $\epsilon_g$, which were as follows:

\[
\begin{align*}
\epsilon_s + \epsilon_d &= 3.85\epsilon_g \\
\epsilon_c + \epsilon_o &= 1.01\epsilon_g \\
\Rightarrow \epsilon_s &= 3.45\epsilon_g \\
\epsilon_o &= 0.61\epsilon_g
\end{align*}
\]

IV. RESULTS AND DISCUSSIONS

Figures 2, 3, and 4 show the relationship between $P_L$ and $m$ for the case $\epsilon_a = 0.0, 0.3, 0.5$, respectively, for gate error values $\epsilon_g = 5.0 \times 10^{-5}, 8.0 \times 10^{-5}, 1.0 \times 10^{-4}, 3.0 \times 10^{-4}$. Note that $P_L$ is the cumulative error for 1000 gates (plus QEC operations with some frequency), and hence may be larger than the underlying error for an individual physical gate, even when operating below threshold. The black circular points are the values for $P_L$ determined from the numerical simulation (section II) and the the blue triangular points are for $P_L$ which are determined by using the exact formula.

Our primary observations are that there is generally good agreement between the data generated by the formula and the simulation (and reasonably good agreement even given the assumption of large $B$, especially with respect to the location of $m_{\text{min}}$), and that $m_{\text{min}}$ is insensitive to variations in $\epsilon_o$ over the range of $\epsilon_a$ considered, with $m_{\text{min}} = 5$ in all cases. Figure 5 shows the behavior of $P_L$ as a function of $\epsilon_a$ for $\epsilon_g = 5 \times 10^{-5}$ and $m = 5$, showing that the variation in $P_L$ is relatively small over lower values of $\epsilon_a$. Figure 6 shows the variation of the $m_{\text{min}}$, with $\epsilon_a$ for $\epsilon_g = 5 \times 10^{-5}$, as given by equation 11.

As expected, within the region $m < m_{\text{min}}$, the error rate is reduced both by increasing $m$ and by increasing $\epsilon_a$, since both result in fewer QEC operations being performed (the only difference being whether the skipped operations are regularly spaced or not), and QEC operations in this region produce more errors, on average, than they correct. Similarly the behavior is reversed for $m > m_{\text{min}}$.

Overall agreement is good; there is a slightly larger disparity for $\epsilon_a = 0$ and small $m$. This is where the largest number of QEC operations occur (since they are attempted frequently and none are skipped), indicating that the approximations in modeling QEC errors are leading to an underestimate of the overall logical error. At higher values of $m$ (where gate errors are a more dominant source of error) agreement is generally better, except at the largest $m$ is large and the approximation is no longer valid.
FIG. 2. $P_L$ vs $m$ for $\varepsilon_g = 1.0 \times 10^{-4}$, $3.0 \times 10^{-4}$, $8.0 \times 10^{-5}$, $5.0 \times 10^{-5}$, where $\varepsilon_a = 0.0$. The triangular blue points are the numerical values of $P_L$ as given by the equation (6). The circular black points are the numerical simulation values of $P_L$.

FIG. 3. $P_L$ vs $m$ for $\varepsilon_g = 5.0 \times 10^{-5}$, $8.0 \times 10^{-5}$, $1.0 \times 10^{-4}$, $3.0 \times 10^{-4}$, where $\varepsilon_a = 0.3$.

V. CONCLUSION

We have found the optimal number of quantum gates to perform before applying an error correction operation. This was done with a semi-abstract model to enable some generality. An analytic expression was presented which provides explicit dependence on the error correction frequency as a function of the gate error rate, ancilla failure rate, and error rates for the correction operation. The various rates depend on the underlying physical gate error rate. The dependence is different for different circuits which are determined by the code used. To be explicit we showed in detail how this works by example. Our example is the commonly used Steane $[[7,1,3]]$ code and
The simulation

The analytical expression

$$\varepsilon_g = 1.0 \times 10^{-4}$$

$$0.050$$

$$0.100$$

$$0.150$$

$$0.200$$

$$0.250$$

$$\varepsilon_g = 3.0 \times 10^{-4}$$

$$0.005$$

$$0.010$$

$$0.015$$

$$0.020$$

$$0.025$$

$$0.030$$

$$\varepsilon_g = 8.0 \times 10^{-5}$$

Logical error rate $P_L$ with $\varepsilon_a = 0.5$

$\varepsilon_g = 5.0 \times 10^{-5}$

$0.005$  $0.010$  $0.015$  $0.020$  $0.025$  $0.030$

$0.0$  $0.05  $0.10  $0.15  $0.20  $0.25$

$0.30$  $0.35  $0.40  $0.45  $0.50  $0.55$

$0.60$  $0.65  $0.70  $0.75  $0.80  $0.85$

$0.90$  $0.95  $1.00$

FIG. 4. $P_L$ vs $m$ for $\varepsilon_g = 5.0 \times 10^{-5}, 8.0 \times 10^{-5}, 1.0 \times 10^{-4}, 3.0 \times 10^{-4}$, where $\varepsilon_a = 0.5$.

Steane ancilla technique. However, we believe an immediate extension could be made to other distance-3 CSS codes, and a more general formula could be derived by a similar technique for other codes.

We have assumed a transversal gate model. Single-qubit logical operations may, however, depending on the computation and code, be dominated by non-transversal gates (such as the $T$ gate). Such gates, which often require preparation of a post-selected ancilla, would require an error model similar to that used for QEC operations. Our treatment is expected to work very well for storage where examples include gates commonly called HOLD, or WAIT.

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TABLE II. Table of possible second-order errors spanning 2 blocks. Such errors can occur (even without skipped QECs) when a QEC error in one block combines with an error (from logic gates or the QEC) in the subsequent block. In B blocks there are \(B - 1\) sets of adjacent blocks, hence the contribution multiplier of \(B - 1\).

| Error source | Error contribution |
|--------------|--------------------|
| \(Q_i(s, d, o)\) | \(mG_{i+1}(s)\) |
| \(Q_i(c)\) | \((B - 1)(1 - \epsilon_a)(\epsilon_s + \epsilon_o)me_g\) |

\[1x\]

\[1x, 1\]

\[1x, 1\]

\[1x, 1\]

\[1x, 1\]

\[1x, 1\]

\[1x, 1\]

\[1x, 1\]

TABLE III. Table of possible second-order errors due to \(f\) skipped QECs spanning \(f + 1\) blocks. This combines errors from operations separated by one or more blocks with skipped QECs, where QECs from blocks \(i\) to \(i + f - 1\) are skipped, and no logical gate errors occur in blocks \(i + 1\) to \(i + f - 1\) (these operations, omitted from the table, are denoted by the column of Xs). The contribution includes contributions for all \(f > 0\), hence the multiplier \(\gamma\), as discussed in section II B.

| Error source | Error contribution |
|--------------|--------------------|
| \(mG_i\) | \(X\) |
| \(mG_{i+1}(s, d, o)\) | \(Q_{i+1}(s)\) |
| \((\text{overall factor 42})\) | \(\gamma(me_g)^2\) |

We found excellent agreement showing that our course-graining (due to a rough classification of error types) provides enough precision to provide a very reliable estimate. Furthermore, we find that the optimum frequency to apply QEC operations is relatively insensitive to ancilla failure probability, over the range of gate error and ancilla error probabilities considered (with the optimum varying from \(m = 3\) to \(m = 6\) but frequency changes within this range making only small differences to the overall logical error), indicating that skipping QEC operations under ancilla failure will in many cases be a successful approach even in a design where QECs are performed infrequently. This will save resources while providing a better overall logical error rate for quantum error correcting codes.

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TABLE IV. Table of possible second-order errors due to $f$ skipped QECs, spanning $f+2$ blocks. As with the analogous Tables I and II, error within QEC operations can lead to a total span of one additional block, compared to the case of Table III, leading to a multiplier $\gamma_3$, as discussed in section IV.B. QECs from blocks $i+1$ to $i+f$ are skipped, and no errors occur in blocks $i+2$ to $i+f$.

| Error source | Error contribution |
|--------------|--------------------|
| $Q_i(s,d,o)$ | $Q_i(c)$ $mG_{i+1+f}$ $Q_{i+1+f}(s)$ (overall factor 42) |
| 1 $s,o$ | $\gamma_3(1-\epsilon_a)m(\epsilon_s + \epsilon_o)\epsilon_g$ |
| 1 $s,o$ | $\gamma_3(1-\epsilon_a)^2(\epsilon_s + \epsilon_o)(\epsilon_s + \epsilon_d)$ |
| 1 | $\gamma_3(1-\epsilon_a)\epsilon_c\epsilon_g$ |
| 1 $s,d$ | $\gamma_3(1-\epsilon_a)^2\epsilon_c(\epsilon_s + \epsilon_d)$ |