Size Distribution of Small Jupiter Trojans in the L₅ Swarm *

KOTOMI UEHATA,₁,₂ TSUYOSHI TERAI,₃ KEIJI OHTSUKI,₁ and FUMI YOSHIDA₄,₅

₁Department of Planetology, Kobe University, Kobe 657-8501, Japan
₂NEC Corporation, Tokyo 108-8001, Japan
₃Subaru Telescope, National Astronomical Observatory of Japan, National Institutes of Natural Sciences, Hilo, HI 96720, USA
₄Planetary Exploration Research Center, Chiba Institute of Technology, Narashino 275-0016, Japan
₅School of Medicine, University of Occupational and Environmental Health, Kita-Kyushu, Japan

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ABSTRACT

We present an analysis of survey observations of the trailing L₅ Jupiter Trojan swarm using the wide-field Hyper Suprime-Cam CCD camera on the 8.2 m Subaru Telescope. We detected 189 L₅ Trojans from our survey that covered about 15 deg² of sky with a detection limit of \( m_r = 24.1 \) mag, and selected an unbiased sample consisting of 87 objects with absolute magnitude \( 14 \leq H_r \leq 17 \) corresponding to diameter \( 2 \text{ km} \leq D \leq 10 \text{ km} \) for analysis of size distribution. We fit their differential magnitude distribution to a single-slope power-law with an index \( \alpha = 0.37 \pm 0.01 \), which corresponds to a cumulative size distribution with an index of \( b = 1.85 \pm 0.05 \). Combining our results with data for known asteroids, we obtained the size distribution of L₅ Jupiter Trojans over the entire size range for \( 9 \leq H_V \leq 17 \), and found that the size distributions of the L₄ and L₅ swarms agree well with each other for a wide range of sizes. This is consistent with the scenario that asteroids in the two swarms originated from the same primordial population. Based on the above results, the ratio of the total number of asteroids with \( D \geq 2 \text{ km} \) in the two swarms is estimated to be \( N_{L4}/N_{L5} = 1.40 \pm 0.15 \), and the total number of L₅ Jupiter Trojans with \( D \geq 1 \text{ km} \) is estimated to be \( 1.1 \times 10^5 \) by extrapolating the obtained distribution.

Keywords: asteroids: general — methods: observational — techniques: photometric

1. INTRODUCTION

The Jupiter Trojans are a group of asteroids that share their orbits with Jupiter and are confined to two swarms centered about the L₄ and L₅ Lagrangian points, which lead and trail Jupiter by about 60 degrees. Although the gravitational effects of the giant planets likely have reduced the number of the swarm asteroids over time, most of Jupiter Trojans have stability timescales comparable to or longer than the age of the Solar System.
System (Levison et al. 1997; Di Sisto et al. 2014). While in situ formation of Jupiter and the Trojans at 5.2 au cannot explain their characteristics such as broad inclination distribution, recent models showed that capture of icy planetesimals from the trans-Neptunian region at the time of orbital instability of giant planets can explain their observed orbital characteristics and estimated total mass (Morbidelli et al. 2005; Nesvorný et al. 2013). On the other hand, a more recent model of capture of Trojans asteroids due to the growth and inward migration of Jupiter shows a higher capture efficiency than the above model based on orbital instability, and can naturally explain the observation that the L₄ swarm is more populated than the L₅ swarm (Pirani et al. 2019). In the latter model, Trojans are originated from the radial location where Jupiter’s core formed. Thus, Trojan asteroids are expected to provide us with clues to understanding origin and radial transport of small icy bodies caused by orbital evolution of giant planets as well as building blocks of the planets (see a review by Emery et al. 2015).

Trojan asteroids are expected to hold unique information about primitive material in the outer Solar System region, and detailed information for some of these asteroids is expected to be revealed by NASA’s Lucy mission (Levison et al. 2021). On the other hand, in order to collect information from a large number of samples for statistical investigation, various observations using ground-based or space telescopes have been carried out (Emery et al. 2015). Here we focus on observation of size distribution of Jupiter Trojans, which can provide us with clues to their origin and collisional evolution. Size distribution of Trojan asteroids has been investigated by several surveys (e.g. Jewitt et al. 2000; Szabó et al. 2007). The slopes of the size distributions of hot Kuiper-belt objects (KBOs) and Jupiter Trojans near their large size end (with absolute magnitudes $H \lesssim 8$, corresponding to diameters $D \gtrsim 100$ km) are shown to be similar but is distinctly shallower than that of cold KBOs (Fraser et al. 2014). This can be interpreted as a support for the view that Jupiter Trojans were captured from scattered KBOs (Fraser et al. 2014), but information for smaller objects will be useful to derive stronger constraints. Size distribution of small Trojan asteroids down to ~2 km has also been examined using the Suprime-Cam CCD camera on Subaru Telescope for the L₄ (Yoshida & Nakamura 2005; Wong and Brown 2015) and L₅ swarms (Yoshida & Nakamura 2008), respectively. It is also known that Jupiter Trojans have two groups with respect to their color or albedo (e.g. Szabó et al. 2007; Emery et al. 2011; Grav et al. 2012), and each group of asteroids are shown to have distinct size distribution (Wong et al. 2014; Wong and Brown 2015). However, in the case of small Trojans in the L₅ swarm, the number of objects used for the analysis of size distribution was limited (62 in Yoshida & Nakamura (2008)) and, furthermore, removal of observational biases may have not been sufficient (§4.1).

On the other hand, a new CCD camera for Subaru Telescope with a larger field of view, called Hyper Suprime-Cam (HSC), has been developed in 2012 (Miyazaki et al. 2012) and became a general user instrument in 2014. Using HSC, we have carried out observational studies of size distributions of L₄ Jupiter Trojans (Yoshida & Terai 2017) and Hilda asteroids (Terai & Yoshida 2018), colors of KBOs (Terai et al. 2018) and Centaurs (Sakugawa et al. 2018), color and size distribution of main-belt asteroids (Maeda et al. 2021), and comparison of size distributions of small bodies in various populations (Yoshida et al. 2019, see also Yoshida et al. (2020) for corrections).

As an extension of this series of works, in the present work, we examine size distribution of small L₅ Jupiter Trojans with HSC. By comparing our results with Yoshida & Terai (2017), who examined size distribution of small L₄ Jupiter Trojans also using HSC, we will discuss implications for the origin of Jupiter Trojans. The organization of this paper is as follows. §2 describes our observation and data analysis, and results are presented in §3. Discussion is given in §4, where we compare our results with previous studies on size distribution of Jupiter Trojans in detail. Our conclusion is given in §5.
Figure 1. Locations of the nine observed fields of our survey. The size of each circle corresponds to the field of view of Hyper Suprime-Cam. Lines of ecliptic latitude ($\beta$) are also shown, with the ecliptic plane shown with the thick line. The background image is taken from Pan-STARRS1 (Chambers et al. 2016).

2. OBSERVATION AND DATA ANALYSIS

Observation of the L5 Trojan swarm of Jupiter was carried out on January 10, 2016 (UT) using the HSC installed at the prime focus of the 8.2 m Subaru Telescope atop Maunakea, Hawaii. The HSC is a gigantic mosaic camera consisting of 104 CCD chips that can be used for scientific observation with a field of view of 1.5 deg in diameter. We surveyed about 15 deg$^2$ of sky in the L5 region near opposition (Figure 1). The right ascensions (R.A.) of the L5 Lagrangian point and opposition at our observation were 103 deg and 111 deg, respectively, and the surveyed area was near the L5 point and was within about 8 deg from opposition. Nine fields of view of HSC shown in Figure 1 were observed in the r-band, and the exposure time for each field was 240 seconds. Each field of view was visited three times with a time interval of $\sim$23 minutes. Observational sequence is shown in Table 1. The average seeing size of each field was between 0$''$.82 and 1$''$.19, allowing us observations under a rather stable condition.

Data analysis was performed following procedures similar to those described in Yoshida & Terai (2017) and Terai & Yoshida (2018). First, the data are processed with the HSC data reduction/analysis pipeline hscPipe (version 3.8.5; Bosch et al. (2018)). We extracted moving object candidates from the source catalogs created by hscPipe. Then we identify L5 Jupiter Trojans based on their apparent motion (see Yoshida & Terai 2017, for details of the procedures). Figure 2 shows the velocities along the ecliptic longitude and latitude of the detected Jupiter Trojan (red) and Hilda (green) candidates. We searched candidate objects for the two populations within the outer boundary shown by the left vertical line and the right curve, and the middle curve was used to distinguish between Hildas and Trojans. We can see that the two populations are clearly separated from each other. From these procedures, we identified 189 L5 Jupiter Trojans and 76 Hildas from our sample. We analyzed the data for the Jupiter Trojans in the present work.

For each of the detected Trojans, we carried out photometric measurements as follows (see Yoshida & Terai (2017) and Terai & Yoshida (2018) for details). First, we determined the center position for each
Table 1. Observation Log

| Field ID | UTC  | Time interval (min: sec) | R.A. (deg) | Decl. (deg) | Airmass | Seeing (arcsec) |
|----------|------|--------------------------|------------|------------|---------|----------------|
| FIELD01  | 06:35:33 | – | 114.14959 | 20.80001 | 1.859 | 1.23 |
|          | 06:58:19 | 22:46 | 114.14958 | 20.80000 | 1.633 | 1.13 |
|          | 07:21:01 | 22:42 | 114.14960 | 20.80000 | 1.466 | 1.20 |
| FIELD02  | 06:40:06 | – | 114.14959 | 22.30001 | 1.796 | 1.19 |
|          | 07:02:51 | 22:45 | 114.14959 | 22.30003 | 1.588 | 1.05 |
|          | 07:25:35 | 22:44 | 114.14959 | 22.30001 | 1.433 | 1.23 |
| FIELD03  | 06:44:38 | – | 112.99960 | 23.40003 | 1.698 | 1.27 |
|          | 07:07:23 | 22:45 | 112.99959 | 23.40001 | 1.516 | 1.04 |
|          | 07:30:08 | 22:45 | 112.99961 | 23.40003 | 1.379 | 1.19 |
| FIELD04  | 06:53:45 | – | 115.54960 | 21.54999 | 1.720 | 1.20 |
|          | 07:16:28 | 22:43 | 115.54960 | 21.54999 | 1.531 | 1.15 |
|          | 07:39:14 | 22:46 | 115.54958 | 21.55000 | 1.389 | 1.14 |
| FIELD05  | 08:12:45 | – | 115.54958 | 20.05001 | 1.241 | 0.90 |
|          | 08:35:35 | 22:50 | 115.54960 | 20.05000 | 1.167 | 0.81 |
|          | 08:58:24 | 22:49 | 115.54960 | 20.05000 | 1.109 | 0.80 |
| FIELD06  | 08:17:19 | – | 116.94958 | 20.80002 | 1.244 | 0.84 |
|          | 08:40:09 | 22:50 | 116.94958 | 20.80002 | 1.169 | 0.83 |
|          | 09:03:03 | 22:54 | 116.94960 | 20.80002 | 1.111 | 0.79 |
| FIELD07  | 08:21:53 | – | 116.94958 | 22.30004 | 1.226 | 0.86 |
|          | 08:44:42 | 22:49 | 116.94960 | 22.30002 | 1.155 | 0.87 |
|          | 09:07:35 | 22:43 | 116.94960 | 22.30004 | 1.101 | 0.75 |
| FIELD08  | 08:26:26 | – | 118.34958 | 21.55001 | 1.230 | 0.90 |
|          | 08:49:18 | 22:52 | 118.34960 | 21.54999 | 1.158 | 0.91 |
|          | 09:12:14 | 22:56 | 118.34959 | 21.55003 | 1.103 | 0.75 |
| FIELD09  | 08:31:00 | – | 118.34960 | 20.04999 | 1.216 | 0.91 |
|          | 08:53:51 | 22:51 | 118.34960 | 20.05001 | 1.147 | 1.00 |
|          | 09:16:47 | 22:56 | 118.34958 | 20.05000 | 1.094 | 0.88 |

of the detected objects using $\chi^2$-fitting of an object model to the image data; the model was generated by integration of Gaussian profiles with the center shifted with a constant motion corresponding to the measured velocity. Then we measure the total flux for each object by aperture photometry (Yoshida & Terai 2017; Terai & Yoshida 2018). The measured flux was converted into apparent magnitude in the AB magnitude system using the photometric zero point, which was estimated in hscPipe based on the data of the Panoramic Survey Telescope and Rapid Response System (Pan-STARRS) 1 survey (Schlafly et al. 2012; Tonry et al. 2012; Magnier et al. 2013). Photometric measurements were performed for 187 of the detected objects, excluding two objects for which the measurement failed due to either a background bright star or a defect in part of the images.

The short observation arcs of the detected objects did not allow us to determine their orbits accurately. However, because the surveyed area was close to opposition, assuming that their orbits are circular we
Figure 2. Distribution of apparent velocities of detected Jupiter Trojans (red dots) and Hildas (green dots) along ecliptic longitude and latitude. The background map shows the distribution of apparent velocities of synthetic Trojans and Hildas, respectively, generated by a Monte Carlo method based on the orbital distribution of known asteroids in each population. The three dashed lines show boundaries used to distinguish detected objects for each population.

were able to estimate the heliocentric distance and orbital inclination of each of the detected objects from its apparent motion. We estimated the heliocentric distance and orbital inclination following Terai et al. (2013, see also Jedicke (1996)). In this case, systematic deviation from the correct value is expected in the estimated heliocentric distance because of the assumption of circular orbits. This systematic errors in estimated heliocentric distance can be corrected as follows (Yoshida & Terai 2017; Terai & Yoshida 2018). First, we generate synthetic orbits of Jupiter Trojans from probability distributions based on the orbital distributions of known Trojans obtained from Minor Planet Center (MPC) database. Then their apparent motions with the same conditions as our observations are calculated, and their heliocentric distances are estimated from their apparent motions in the manner described above. The left panel of Figure 3 shows a comparison between their estimated heliocentric distance ($R_{est}$) and the true generated values ($R_{true}$), where the systematic deviation between the two values can be seen. The solid line represents the best-fit linear function to the obtained distribution given as

$$R_{est} = 0.497R_{true} + 2.596.$$  

Using this derived relationship, we re-estimated the heliocentric distance of all the synthetic objects and confirmed that their true values can be recovered (Figure 3, right panel). The statistical error of the corrected

1 http://minorplanetcenter.net/iau/lists/JupiterTrojans.html
Figure 3. Left: Relationship between true generated heliocentric distance of synthetic Trojans and the values of the heliocentric distance estimated from their apparent motion under the same observation conditions as our survey. The systematic deviation between the two values is due to the assumption of circular orbits in the estimation from apparent motions. The solid line represents the best-fit to the estimated values given by Eq.(1). Right: Relationship between the generated heliocentric distance and re-estimated values. The latter was obtained by removing the systematic deviation using Eq.(1).

The heliocentric distance was about 0.09 au (Yoshida & Terai 2017). Figure 4 shows the relationship between the derived heliocentric distance and orbital inclination together with their histograms for our 189 Trojan asteroids. Random errors for the obtained heliocentric distances and orbital inclinations are 0.18 au and 1.5 deg, respectively.

Detectability of small objects such as Trojan asteroids depends on their apparent magnitudes as well as image quality. Therefore, for an accurate investigation of these bodies, evaluation of detection completeness is important for removal of such biases. In the case of HSC, image quality as well as effective area for detection of moving objects differs for each CCD, and the detection efficiency also depends on the observation condition for each visit of the field. In the present work, we examined the detection efficiency by implanting synthetic moving objects into the image obtained for each field and for each visit, and processing them with hscPipe in the same manner as for the actual detected moving objects. Detailed procedures are similar to those described in Yoshida & Terai (2017, their §3.3 and Figure 5), but here we improved the method by implanting synthetic objects into actual processed images rather than synthetic blank images as done in Yoshida & Terai (2017). This improvement allowed us more accurate estimate of the detection efficiency for each image used for our size distribution analysis. The obtained detection efficiency as a function of apparent magnitude $m$ is then fit by the following function (Yoshida & Terai 2017):

$$\eta(m) = \eta_0 \sum_{k=1,2} \frac{\epsilon_k(m)}{2} \left[ 1 - \tanh \left( \frac{m - m_{50}}{w_k} \right) \right]$$

(2)
**Figure 4.** Relationship between the derived heliocentric distances and orbital inclinations together with their histograms for the detected 189 Trojan asteroids. Random errors for the obtained heliocentric distances and orbital inclinations are 0.18 au and 1.5 deg, respectively.

where

\[ \epsilon_1(m) = 1 - \epsilon_2(m) = \frac{1}{2} \left[ 1 - \tanh \left( \frac{m - m_{50}}{0.2} \right) \right]. \] (3)

In the above, \( \eta_0 \) is the maximum detection efficiency, \( m_{50} \) is the apparent magnitude where the efficiency is \( \eta_0/2 \), and \( w_k \) (\( k = 1, 2 \)) are the transition widths; all of these are evaluated by fitting of the calculated detection efficiencies of the implanted synthetic objects. We examined the detection efficiency for CCDs representative of different locations of each field (i.e., near the center, near the edge, and at an intermediate region), and found that the detection efficiency for almost all the CCDs exceeds 0.5 (i.e., more than 50% detection) when the apparent magnitude of the implanted synthetic object is 24.1 mag. Thus, we set 24.1 mag as the detection limit of our survey.
3. RESULTS

3.1. Sample Selection

In order to derive size distribution of detected Trojans, the measured apparent magnitude for each object was converted into the absolute magnitude in the $r$-band as (Bowell et al. 1989)

$$H_r = m_r - 5 \log(R\Delta) - P(\theta),$$

(4)

where $m_r$ is the apparent magnitude in the $r$-band; $R$ and $\Delta$ are the heliocentric and geocentric distances, respectively; and $P(\theta)$ is the phase function for a solar phase angle $\theta$ given by

$$P(\theta) = -2.5 \log \left[ (1 - G)\Phi_1 + G\Phi_2 \right].$$

(5)

Here, $G$ is the slope parameter, and $\Phi_1$ and $\Phi_2$ are the phase functions given by

$$\Phi_i = \exp(-A_i[\tan(\theta/2)]^{B_i}) \quad (i = 1, 2)$$

(6)

with $A_1 = 3.33$, $A_2 = 1.87$, $B_1 = 0.63$, and $B_2 = 1.22$. Because the observations were carried out only for one night, we were not able to determine $P(\theta)$ from our observations. Thus, following previous works (e.g. Grav et al. 2011; Yoshida & Terai 2017), we assumed $G = 0.15$, which is a typical value for main-belt asteroids.

Figure 5 shows the plots of the absolute magnitudes as a function of the heliocentric distance for each of the 187 Trojans for which photometric measurements were performed. The uncertainty in the absolute magnitude comes from errors in the estimates of the apparent magnitude as well as the heliocentric and geocentric distances. In our observations, measurement errors in apparent magnitude were less than $\sim 0.05$ mag for objects with $m_r \leq 23$ mag, while they were as large as 0.15 mag for $m_r \approx 24$ mag. Also, errors in the estimated heliocentric and geocentric distances were about 0.09 au (Yoshida & Terai 2017). Taking all these into account, errors in the obtained absolute magnitude of each object are shown with error bars in Figure 5.

The dashed line in Figure 5 represents the apparent magnitude of the detection limit, $m_r = 24.1$ mag, derived in §2. In order to remove detection bias caused by the decreasing observed brightness with increasing distance from the Sun and Earth, we defined the outer edge of the heliocentric distance of detected Trojans by $R = 5.5$ au, where $m_r = 24.1$ mag corresponds to $H_r = 17.0$ mag, and selected objects that satisfy $R \leq 5.5$ au and $H_r \leq 17.0$ mag as our unbiased sample. The extracted sample consists of 87 Jupiter Trojans.

3.2. Size Distribution

The obtained absolute magnitude $H_r$ for each object can be converted into its diameter $D$ by (Pravec & Harris 2007)

$$\log D = 0.2m_{\odot,r} + \log(2R_{\oplus}) - 0.5 \log p - 0.2H_r,$$

(7)

where $m_{\odot,r}$ is the apparent $r$-band magnitude of the Sun ($-26.91$ mag; Fukugita et al. (2011)), $R_{\oplus}$ is the heliocentric distance of Earth (i.e. 1au) in the same units as $D$, and $p$ is the geometric albedo. We assumed $p = 0.05$ on the basis of a recent analysis of Near-Earth Object Wide-field Infrared Survey Explorer (NEOWISE) data for Jupiter Trojans with $D \geq 10$ km (Romanishin & Tegler 2018).
Figure 5. Absolute magnitudes with error bars as a function of heliocentric distance for the 187 Trojan asteroids for which photometric measurements were successfully performed. The dashed lines shows the limiting apparent magnitude of 24.1 mag. The solid lines represent the boundaries we set to select our unbiased sample, i.e., $R = 5.5$ au and $H_r = 17.0$.

We derived absolute magnitude distribution from our unbiased sample consisting of 87 objects, which can be converted into their size distribution. Here, the derived number of objects should be corrected with the detection efficiency. Suppose that the detection efficiency at $i$-th visit ($i = 1, 2, 3$ in our observations) for an object $j$ with its apparent magnitude $m_j$ is given by $\eta_i(m_j)$. Then the efficiency for the detection of this object in all the three visits can be given as

$$\eta(m_j) = \prod_{i=1}^{3} \eta_i(m_j).$$

Thus, the corrected cumulative number of objects with absolute magnitude smaller than $H_r$ is

$$N(< H_r) = \sum_{j: H_j < H_r} \frac{1}{\eta(m_j)}.$$  

Figure 6 shows the plots of the cumulative absolute magnitude distribution of the $L_5$ Jupiter Trojans in our unbiased sample of 87 objects with $H_r < 17.0$ mag. The corresponding diameter values are also indicated on the upper horizontal axis, which show that the Trojan asteroids in our sample have $D \simeq 2 - 10$ km. We performed fitting of the distribution by a power-law with a single slope as follows. The differential absolute magnitude distribution function defined by $\Sigma(H) = dN(< H_r)/dH$ for a single-slope power-law can be written as
Figure 6. Cumulative absolute magnitude distribution for the 87 unbiased L₄ Jupiter Trojans obtained from our survey (red dots with error bars). Note that the cumulative numbers are corrected using the detection efficiency as shown in Eq.(9). Corresponding diameters for the assumed albedo of 0.05 (Romanishin & Tegler 2018) are also shown at the upper horizontal axis. Best-fit single-slope power-law distribution is shown with the dashed line.

\[
\Sigma(H) = 10^{\alpha(H-H_0)},
\]

(10)

where \(\alpha\) is the power-law index and \(H_0\) is defined so that \(\Sigma(H_0) = 1\). For a constant albedo, the cumulative diameter distribution can be written as

\[
N(> D) \propto D^{-b}
\]

(11)

with \(b = 5\alpha\).

We used the maximum-likelihood method (e.g. Bernstein et al. 2004) for fitting Eq.(10) to the obtained \(H_r\) distribution. For the numerical analysis of the maximum-likelihood estimation, we used the Markov Chain Monte Carlo method (MCMC) with a Python package emcee

² (Foreman-Mackey et al. 2013). Figure 7 shows the post-event distributions of sampling for \(\alpha\) and \(H_0\) by MCMC. We obtained the best-fit power-law index of \(\alpha = 0.37 \pm 0.01\) (which corresponds to \(b = 1.85 \pm 0.05\)) with \(H_0 = 11.55 \pm 0.15\), which is shown in Figure 6 with the dashed line. This obtained slope agrees with the result of Yoshida & Terai (2017), who also used data obtained using HSC and found \(\alpha = 0.37 \pm 0.01\) for 431 small L₄ Trojans in a similar size range. This shows that the slope of the size distribution of L₄ and L₅ Trojans agree with each other for \(D \approx 2 - 10\) km.

² https://emcee.readthedocs.io/en/stable/
Next, we compare size distributions of L$_4$ and L$_5$ Trojans for a wider range of diameters, by combining the data obtained by HSC (Yoshida & Terai 2017 and this work) with those for known Jupiter Trojans. The left panel of Figure 8 shows the distributions of V-band absolute magnitude and estimated diameters for L$_4$ Trojans; the green circles show the data obtained by HSC (Yoshida & Terai 2017) and the green squares show the data for known Jupiter Trojans obtained from the MPC catalog. In order to compare with the MPC data, which are given in V-band magnitude, we assumed the $V - r$ color to be 0.25 mag (Szabó et al. 2007) and scaled the distribution for the HSC data so that they match the MPC data at $H_V = 14.0$ mag (Yoshida & Terai 2017). Yoshida & Terai (2017) showed that the combined distribution can be fitted by a broken power-law with $\alpha_1 = 0.50 \pm 0.01$ for $H_V \leq H_B$ and $\alpha_2 = 0.37 \pm 0.01$ for $H_V > H_B$ with break magnitude $H_B = 13.56^{+0.04}_{-0.06}$ mag.

The right panel of Figure 8 shows similar plots for L$_5$ Jupiter Trojans, where the HSC data obtained in the present work are combined with the data for known L$_5$ Trojans taken from the MPC catalog. We converted the $r$-band magnitude of our sample into the $V$-band magnitude in the same way as in Yoshida & Terai (2017). Then we scaled the best-fit single-slope power-law distribution to our HSC data (Figure 6) so that it matches the MPC data at $H_V = 14.0$ mag, avoiding the use of the MPC data near small size end where the slope of the magnitude distribution begins to decay. These plots show similarity of the size distributions of the two swarms for a wide range of sizes. We calculated the slope of the magnitude distributions at $11.0 < H_V < 14.0$ using data given in the MPC catalog (updated April 23, 2021) by least-square fitting and found that $\alpha = 0.48 \pm 0.03$ for L$_4$ and $\alpha = 0.46 \pm 0.03$ for L$_5$, respectively. Figure 9 shows the combined
Figure 8. Left: Combined absolute magnitude distribution for the L4 swarm of the known Trojan asteroids taken from the MPC catalog (green squares) and those from the HSC observation by Yoshida & Terai (2017, green circles) scaled at $H_V = 14.0$ mag. The $V - r$ color was assumed to be 0.25 mag (Szabó et al. 2007) in converting the HSC $r$-band magnitudes into the $V$-band magnitudes. Corresponding diameters for the assumed albedo of 0.05 (Romanishin & Tegler 2018) are also shown at the upper horizontal axis. Note that Yoshida & Terai (2017) assumed an albedo of 0.07 based on Grav et al. (2012), while we used 0.05 for both swarms based on the above more recent work. Right: Same as the left panel but for the L5 swarm. The HSC data (red circles) in this case are those obtained by the present work, and are scaled so that their best-fit single-slope power-law distribution (dashed line; see also Figure 6) matches the MPC data (magenta squares) at $H_V = 14.0$ mag (see text).

cumulative magnitude distributions for the L4 and L5 swarms normalized at their values at $H_V = 17.0$ mag. This clearly demonstrates that the size distributions of L4 and L5 Trojans agree well with each other in the size range of $D \geq 2$ km.

By comparing the combined size distribution for the L5 swarm based on our results (right panel of Figure 8) with that for the L4 swarm based on Yoshida & Terai (2017, left panel of Figure 8), we find that the ratio of the total numbers of L4 and L5 Trojan asteroids with $D \geq 2$ km is estimated to be $N_{L4}/N_{L5} = 1.40 \pm 0.15$, which is consistent with previous works that obtained the ratio for larger asteroids (e.g. Grav et al. 2012, see §4.2). Extrapolating the newly obtained size distribution, the total number of L4 Jupiter Trojans with $D \geq 1$ km is estimated to be $1.1 \times 10^5$, while L4 Jupiter Trojans with $D \geq 1$ km is estimated to be $1.5 \times 10^5$ from the extrapolation of the results of Yoshida & Terai (2017). Combining these, the total number of Jupiter Trojans (including both L4 and L5 swarms) with $D \geq 1$ km is estimated to be $2.6 \times 10^5$. On the other hand, using data also obtained by the Subaru HSC, Maeda et al. (2021) show that the total number of the main-belt asteroids with $D \geq 1$ km is about $2 \times 10^6$ (their Figure 10), although there is some uncertainty in the estimate. This
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Figure 9. Combined absolute magnitude distribution for the L_4 (squares) and L_5 (circles) swarms normalized by their values at H_V = 17.0 mag.

indicates that the total number of Jupiter Trojans with D ≥ 1 km is significantly (probably by nearly an order of magnitude) smaller than that of the main-belt asteroids in the same size range.

4. DISCUSSION

4.1. Comparison with Other Previous Studies

We have shown that the results for our L_5 Jupiter Trojan sample are consistent with those obtained by Yoshida & Terai (2017) for the L_4 swarm. Here we compare our results with other previous works on the size distribution of Jupiter Trojans (Table 2). Jewitt et al. (2000) derived size distribution of L_4 Jupiter Trojans from observations using the University of Hawaii 2.2 m telescope, and obtained α = 0.40 ± 0.05 (b = 2.0 ± 0.25) for 11 ≤ H ≤ 16. Szabó et al. (2007) derived size distribution for the whole Jupiter Trojans including both L_4 and L_5 swarms using the SDSS Moving Object Catalog 3 (MOC3), and obtained α = 0.44 ± 0.05 (b = 2.2 ± 0.25) for 9 ≤ H ≤ 13.5. Grav et al. (2011) derived size distributions of L_4 and L_5 Jupiter Trojans using data obtained by the NEOWISE survey, and found that the distributions for the two swarms are similar with α = 0.4 (b = 2) for 20km ≤ D ≤ 80 km. Using absolute magnitudes of Jupiter Trojans reported by the MPC, Fraser et al. (2014) obtained α_1 = 1.0 ± 0.2 and α_2 = 0.36 ± 0.01 for the large and small Jupiter Trojans, respectively, with a break at H = 8.4^{+0.2}_{-0.1}. Wong et al. (2014) examined size distributions of both L_4 and L_5 Jupiter Trojans with H < 12.3 using the SDSS MOC4 data, and found that the two swarm distributions are indistinguishable within uncertainties. Their results showed that α = 0.46 ± 0.01 for H ≥ 8.16. Yoshida & Terai (2017) showed that their results are consistent with
these previous works for the size range that overlaps with their HSC data. Our results are also consistent with the above previous studies. The similarity of the size distributions of the two swarms found by some of these previous works (e.g. Grav et al. 2011; Wong et al. 2014) is also consistent with the present work, which confirmed that similar size distributions are extended to smaller sizes.

Size distributions of small Jupiter Trojans ($D < 10$ km) have been investigated in detail using the Suprime-Cam CCD camera on Subaru Telescope. Using a dataset obtained by a survey primarily targeting at main-belt asteroids (Yoshida et al. 2003), Yoshida & Nakamura (2005) detected 51 L4 Jupiter Trojans and derived their size distribution. They found that the observed size distribution can be fit either by a single slope power-law with $b = 1.89$ ($\alpha = 0.378$) for $14 < H < 17.7$, or by a broken power-law with $b_1 = 2.39$ ($\alpha_1 = 0.48$) for $H < 16$ and $b_2 = 1.28$ ($\alpha_2 = 0.26$) for $H > 16$. The power-law index for the above single-slope power-law distribution is consistent with the present work and Yoshida & Terai (2017). Using a dataset obtained by another survey also primarily targeting at main-belt asteroids (Yoshida & Nakamura 2007), Yoshida & Nakamura (2008) detected 62 L5 Jupiter Trojans and derived their size distribution. Their results showed that $b = 2.1 \pm 0.3$ ($\alpha = 0.42 \pm 0.06$) for $15.5 \leq H \leq 17.5$, but they did not find a break at $H \approx 16$, which was found by Yoshida & Nakamura (2005). Using HSC, which has a field of view about seven times as large as that of Suprime-Cam, Yoshida & Terai (2017) and the present work performed surveys for the area of sky of about 26 deg$^2$ and about 15 deg$^2$, respectively; detected 631 and 189 Jupiter Trojans, respectively; and derived size distribution using unbiased samples consisting of 481 and 87 Trojans, respectively. On the other hand, using Suprime-Cam, Yoshida & Nakamura (2005) and Yoshida & Nakamura (2008) carried out surveys for about 3 deg$^2$ and about 4 deg$^2$, respectively; detected 51 and 62 Jupiter Trojans, respectively; and derived size distribution using all these detected objects, but removal of observational bias may have not been sufficient. Thus, it is possible that the difference in the size distributions found by these previous studies near the small size end such as the existence of a break at $H \approx 16$ may be due to either the small number of detected objects and/or observational bias. More recently, Wong and Brown (2015) carried out a survey of L4 Jupiter Trojans using Suprime-Cam, and obtained $b_0 \approx 4.55$ ($\alpha_0 \approx 0.91$) for $H < 8.46$, $b_1 \approx 2.2$ ($\alpha_1 \approx 0.44$) for $8.46 < H < 14.9$, and $b_2 \approx 1.8$ ($\alpha_2 \approx 0.36$) for $H > 14.9$. These results are roughly consistent with the present work and Yoshida & Terai (2017). All these works show that the size distributions of small Jupiter Trojans are similar for the L4 and L5 swarms, and can be approximated by $b \approx 4.5 - 5$ ($\alpha \approx 0.9 - 1.0$) for $H \leq 8.5 - 9$; $b \approx 2.0 - 2.5$ ($\alpha \approx 0.40 - 0.50$) for $9 \leq H < H_B$; and $b \approx 1.8 - 1.9$ ($\alpha \approx 0.36 - 0.38$) for $H \geq H_B$, with $H_B \approx 13.5 - 15$ (corresponding break diameter $\approx 5 - 10$ km).

On the other hand, it has been shown that the spectral slopes and albedos of Jupiter Trojans can be divided into two groups, a "red" group consistent with the asteroidal D-type and a "less-red" group consistent with the asteroidal P-type (e.g. Szabo et al. 2007; Emery et al. 2011; Grav et al. 2012; Wong et al. 2014). While the red-sloped spectra observed at visible and near-infrared wavelengths are often interpreted as indicating the existence of organic material, absence of strong absorptions expected for such material has been reported and the observed spectra may be explained by amorphous and/or space-weathered silicates (Emery & Brown 2003, 2004; Emery et al. 2011, 2015). Wong et al. (2014) examined the magnitude distributions of Jupiter Trojans in each of these two color populations with $H < 12.3$ and found that both the bright-end and the faint-end slopes are different (i.e., $\alpha_{\text{R}}^b = 0.97^{+0.05}_{-0.04}$ versus $\alpha_{\text{LR}}^b = 1.25^{+0.09}_{-0.04}$ for the bright-end and $\alpha_{\text{R}}^f = 0.38 \pm 0.02$ versus $\alpha_{\text{LR}}^f = 0.52^{+0.03}_{-0.01}$ for the faint-end, where "R" and "LR" denote red and less-red color populations, respectively). Wong and Brown (2015) performed observations of colors and magnitude distributions of 557 small L4 Jupiter Trojans using Suprime-Cam on Subaru Telescope with $g'$- and $i'$-
filters. In this case, most likely owing to the effect of asteroid rotation, the bimodality in color distribution that was seen for larger objects was not seen. However, a tendency of decreasing $g-i$ color with increasing magnitude was confirmed, which is consistent with the difference in the slopes of magnitude distributions found by Wong et al. (2014). Such a difference in size distributions for the two color populations may reflect origins and collisional evolution of these bodies (Wong and Brown 2015, 2016), and further observations as well as modeling efforts would be desirable.

4.2. Implications for the Origin of Jupiter Trojans

The results of the present study as well as previous works on the size distribution of Jupiter Trojans show that the size distributions of the $L_4$ and $L_5$ swarms agree well with each other for a wide range of sizes. Fraser et al. (2014) compiled Kuiper Belt survey data and compared the absolute magnitude distributions of hot and cold Kuiper-belt populations with Jupiter Trojans. They found that, for the hot KBOs and Jupiter Trojans, the bright and faint object slopes and the break magnitude (when corrected for the bimodal albedo distribution for the hot KBOs) are statistically indistinguishable, while the cold classical KBOs show a

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**Table 2.** Slope Indices and Break Magnitudes for Jupiter Trojan Absolute Magnitude Distribution

| Reference                    | Swarm | $\alpha_0$ ($H < H_{B0}$) | $\alpha_1$ ($H_{B0} < H < H_B$) | $\alpha_2$ ($H < H_B$) | $H_{B0}, H_B$ |
|------------------------------|-------|-----------------------------|----------------------------------|-------------------------|---------------|
| This work                    | $L_4$ | 0.48 $\pm$ 0.03             |                                  |                         |               |
|                             | $L_5$ | 0.46 $\pm$ 0.03             | 0.37 $\pm$ 0.01                  |                         |               |
| Yoshida & Terai (2017)       | $L_4$ | 0.50 $\pm$ 0.01             | 0.37 $\pm$ 0.01                  |                         | $H_B = 13.56 \pm 0.04$ |
| Wong and Brown (2015)        | $L_4$ | 0.91 $\pm 0.19$ (c) $-0.16$ | 0.45 $\pm 0.05$ (c,d) $-0.16$   | 0.36 $\pm 0.05$ (c) $-0.09$ | $H_{B0} = 8.46 \pm 0.49$ $H_B = 14.93 \pm 0.73$ |
|                             | $L_5$ | 0.84 $\pm 0.22$ (e) $-0.16$ | 0.33 $\pm 0.04$ (e) $-0.03$     |                         |               |
|                             |       | 0.61 $\pm 0.07$ (f,g)       |                                  |                         |               |
| Wong et al. (2014)           | $L_4$, $L_5$ | 1.11 $\pm 0.02$ (c)         | 0.46 $\pm 0.01$ (c)             |                         | $H_{B0} = 8.16 \pm 0.03$ |
|                             |       | 0.97 $\pm 0.05$ (e) $-0.04$ | 0.38 $\pm 0.02$ (e)             |                         | $H_{B0} = 8.70 \pm 0.08$ |
|                             |       | 1.25 $\pm 0.09$ (f) $-0.04$ | 0.52 $\pm 0.03$ (f) $-0.01$     |                         | $H_{B0} = 8.15 \pm 0.06$ |
| Fraser et al. (2014)         | $L_4$, $L_5$ | 1.0 $\pm 0.2$               | 0.36 $\pm 0.01$                 |                         | $H_{B0} = 8.4 \pm 0.2$ |
| Grav et al. (2011)           | $L_4$, $L_5$ | 0.4                          |                                  |                         |               |
| Yoshida & Nakamura (2008)    | $L_5$ |                                  |                                  | 0.42 $\pm 0.06$        |               |
| Szabó et al. (2007)          | $L_4$, $L_5$ |                                  |                                  | 0.44 $\pm 0.05$        |               |
| Yoshida & Nakamura (2005)    | $L_4$ |                                  |                                  | 0.38 $\pm 0.02$ (b)    |               |
| Jewitt et al. (2000)         | $L_4$, $L_5$ | 0.89 $\pm 0.15$             |                                  |                         | $H_{B0} = 10.2 \pm 0.5$ |
|                             | $L_4$ |                                  | 0.40 $\pm 0.05$                 |                         |               |

Note— (a) $H_{B0} \approx 8-9$ (Fraser et al. 2014; Wong et al. 2014; Wong and Brown 2015), although an earlier work suggested a somewhat larger value (Jewitt et al. 2000). (b) $H_B \approx 13.5-15$ (Wong and Brown 2015; Yoshida & Terai 2017). (c) All sample. (d) This was obtained from fitting to their Subaru observation data, while $\alpha_1 = 0.43 \pm 0.02$ was obtained from fitting to bright Trojans in the MPC catalog. (e) Red objects only. (f) Less red objects only. (g) The distribution for the less-red objects with $H < 12.3$ was found to be consistent with a single power law with this slope index. (h) As an alternative distribution, the one with an additional break at $H \approx 16$, $\alpha = 0.48 \pm 0.02$ for $H \leq 16$ and $0.26 \pm 0.02$ for $H \geq 16$, was also suggested.
much steeper large object slope. Fraser et al. (2014) argued that the similarity between the magnitude distributions of the hot KBOs and Jupiter Trojans is consistent with the scenario that they originated from the same primordial population and were scattered to their current locations afterwards (Morbidelli et al. 2005; Nesvorný et al. 2013). The similarity between the size distributions of L₄ and L₅ Jupiter Trojans that we confirmed in the present work is also compatible with the view that these Trojan asteroids originated from the same primordial population.

| Table 3. Degree of L₄/L₅ Asymmetry |
|-----------------------------------|
| Reference                         | \( N_{L4}/N_{L5} \)  |
| This work                         | 1.40 ± 0.15           |
| Grav et al. (2012)                | 1.34                  |
| Grav et al. (2011)                | 1.4 ± 0.2             |
| Nakamura & Yoshida (2008)         | 1.9 ± 0.4             |
| Szabó et al. (2007)               | 1.6 ± 0.1             |

On the other hand, it is known that the L₄ and L₅ swarms have asymmetry in terms of total number (Table 3). Szabó et al. (2007) estimated from the SDSS MOC 3 data that \( N_{L4}/N_{L5} = 1.6 \pm 0.1 \). On the basis of the number of detected Trojans for the two swarms by Subaru Suprime-Cam (Yoshida & Nakamura 2005, 2008) and a model for their surface number density distribution, Nakamura & Yoshida (2008) estimated the total number of Trojans with \( D \geq 2 \) km in each swarm, which leads to \( N_{L4}/N_{L5} = 1.9 \pm 0.4 \). Using the data of Jupiter Trojans with \( D > 10 \) km seen by WISE, Grav et al. (2011) obtained \( N_{L4}/N_{L5} = 1.4 \pm 0.2 \), while Grav et al. (2012) found \( N_{L4}/N_{L5} \approx 1.34 \) also using the WISE data but for large objects with \( D > 50 \) km. Our results show that \( N_{L4}/N_{L5} = 1.40 \pm 0.15 \) for \( D \geq 2 \) km (§3.2), which is consistent with these previous works. The model of capture of Trojans based on the original Nice model (Morbidelli et al. 2005) predicts equal efficiency of capture into the L₄ and L₅ regions and cannot explain the observed asymmetry. On the other hand, the model of capture based on the so-called "jumping Jupiter model" (Nesvorný & Morbidelli 2012) is potentially capable of producing an asymmetry by late passages of an ice giant near the L₅ region that may have depleted the L₅ population (Nesvorný et al. 2013), and the model of capture due to the growth and inward migration of Jupiter can naturally explain the observed asymmetry (Pirani et al. 2019). Although the results of our present study cannot specify mechanisms that caused the asymmetry, our results suggest that the mechanism that lead to the asymmetry did not cause a noticeable difference between the size distributions of the two swarms.

Recently, Hellmich et al. (2019) investigated the influence of the Yarkovsky effect on the long-term orbital evolution of Jupiter Trojans. They found that objects with radii \( R \lesssim 1 \) km are significantly influenced by the Yarkovsky effect and these small bodies can be depleted over timescales of \( \sim 10^8 \) years, creating a shallower slope in the size distribution with a turning point at \( 100 \) m \( \lesssim R \lesssim 1 \) km. This is a size range below that of the currently observable population, thus we cannot test this model with the currently available observational data. Future observations of sub-km Jupiter Trojans would provide further constraints on their origin and evolution.

5. CONCLUSION

In the present work, using data obtained by HSC on Subaru Telescope, we derived size distribution of small Jupiter Trojans in the L₅ swarm with \( 14 \leq H \leq 17 \), corresponding to \( 2 \) km \( \leq D \leq 10 \) km. We
found that the obtained distribution can be approximated by a single-slope power-law, with an index for the
differential magnitude distribution \( \alpha = 0.37 \pm 0.01 \), which corresponds to the cumulative size distribution
for an assumed constant albedo with an index of \( b = 1.85 \pm 0.05 \). This is consistent with Yoshida & Terai
(2017), who examined the size distribution of small Jupiter Trojans in the L4 swarm also using HSC. By
combining our results with the size distribution of known L5 Trojans, we also obtained the size distribution
of L5 Jupiter Trojans over the entire size range for \( 9 \lesssim H_V \lesssim 17 \). We also found that the slopes of the size
distributions of L4 and L5 Trojans agree well with each other at \( H \gtrsim 9 \), and their power-law indices are
approximately given by \( b \approx 2.0 - 2.2 (\alpha \approx 0.40 - 0.44) \) for \( 9 < H < H_B \) and \( b \approx 1.8 - 1.9 (\alpha \approx 0.36 - 0.38) \)
for \( H \gtrsim H_B \) with \( H_B \approx 13.5 - 15 \) (corresponding break diameter \( \approx 5 - 10 \) km). Extrapolating the newly
obtained size distribution, the total number of L5 Jupiter Trojans with \( D \geq 1 \) km is estimated to be \( 1.1 \times 10^5 \).
Combining this estimate with the result for the Jupiter Trojans in the L4 swarm in the same range (Yoshida
& Terai 2017), the total number of Jupiter Trojans with \( D \geq 1 \) km in the L4 and L5 swarms is estimated
\[ N_{L4} = 2.6 \times 10^5, \text{ which is significantly smaller than the recent estimate of the total number of the main-belt}\]
asteroids in the same size range (\( \sim 2 \times 10^6 \); Maeda et al. 2021).

It is known that the L4 and L5 swarms have asymmetry in their total number, with the L4 swarm having
significantly more objects than the L5 swarm. Combining our results for the L5 swarm with those obtained
by Yoshida & Terai (2017) for the L4 swarm, the ratio of the total number of asteroids with \( D \geq 2 \) km in the
two swarms is estimated to be \( \frac{N_{L4}}{N_{L5}} = 1.40 \pm 0.15 \). Our results suggest that the mechanism that caused this
asymmetry did not cause a noticeable difference between the size distributions of the two swarms. Future
observations of still smaller asteroids with larger telescopes and/or detailed close observations of asteroid
surfaces by the Lucy spacecraft are expected to provide further constraints on the origin and evolution of
Jupiter Trojans and the Solar System.

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REFERENCES

Bernstein, G. M., Trilling, D. E., Allen, R. L., et al. 2004, AJ, 128, 1364.
Bosch, J., Armstrong, R., Bickerton, S., et al. 2018, PASJ, 70, S5.
Bowell, E., Hapke, B., Domingue, D., et al. 1989, Asteroids II, 524
Chambers, K. C., Magnier, E. A., Metcalfe, N., et al. 2016, arXiv:1612.05560
Di Sisto, R. P., Ramos, X. S., & Beaugé, C. 2014, Icarus, 243, 287
Emery, J. P. & Brown, R. H. 2003, Icarus, 164, 104
Emery, J. P. & Brown, R. H. 2004, Icarus, 170, 131
Emery, J. P., Burr, D. M., & Cruikshank, D. P. 2011, AJ, 141, 25
Emery, J. P., Marzari, F., Morbidelli, A., French, L. M., & Grav, T. 2015, in Asteroids IV, ed. Michel, P.,
DeMeo, F. E., & Bottke, W. F. (Tucson, University of Arizona Press), 203
Foreman-Mackey, D., Hogg, D. W., Lang, D., et al. 2013, PASP, 125, 306
Fraser, W. C., Brown, M. E., Morbidelli, A., Parker, A., Batygin, K. 2014, ApJ, 782, 100
Fukugita, M., Yasuda, N., Doi, M., et al. 2011, AJ, 141, 47
Grav, T., Mainzer, A. K., Bauer, J., et al. 2011, ApJ, 742, 40
Grav, T., Mainzer, A. K., Bauer, J. M., et al. 2012, ApJ, 759, 49
Hellmich, S., Mottola, S., Hahn, G., et al. 2019, A&A, 630, A148
Jedicke, R. 1996, AJ, 111, 970
Jewitt, D. C., Trujillo, C. A., & Luu, J. X. 2000, AJ, 120, 1140
Levison, H. F., Shoemaker, E. M., & Shoemaker, C. S. 1997, Nature, 385, 42
Levison, H. F., Olkin, C. B., Noll, K. S., et al. 2021, PSJ, 2, 171.
Maeda, N., Terai, T., Ohtsuki, K., et al. 2021, AJ, 162, 280
Magnier, E. A., Schlafly, E., Finkbeiner, D., et al. 2013, ApJS, 205, 20.
Miyazaki, S., et al. 2012, Proc. SPIE, 8446, 84460Z
Morbidelli, A., Levison, H. F., Tsiganis, K., & Gomes, R. 2005, Nature, 435, 462
Nakamura, T. & Yoshida, F. 2008, PASJ, 60, 293
Nesvorný, D. & Morbidelli, A. 2012, AJ, 144, 117
Nesvorný, D., Vokrouhlický, D., & Morbidelli, A. 2013, ApJ, 768, 45
Pirani, S., Johansen, A., Bitsch, B., et al. 2019, A&A, 623, A169
Pravec, P. & Harris, A. W. 2007, Icarus, 190, 250.
Romanishin, W., & Tegler, S. C. 2018, AJ, 156, 19
Sakugawa, H., Terai, T., Ohtsuki, K., et al. 2018, PASJ, 70, 116
Schlafly, E. F., Finkbeiner, D. P., Jurić, M., et al. 2012, ApJ, 756, 158
Szabó, G. M., Ivezic, Ž., Jurić, M., et al. 2007, MNRAS, 377, 1393
Terai, T. & Yoshida, F. 2018, AJ, 156, 30
Terai, T., Takahashi, J., & Itoh, Y. 2013, AJ, 146, 111
Terai, T., Yoshida, F., Ohtsuki, K., et al. 2018, PASJ, 70, S40
Tonry, J. L., Stubbs, C. W., Lykke, K. R., et al. 2012, ApJ, 750, 99
Wong, I., Brown, M. E., & Emery, J. P. 2014, AJ, 148, 112
Wong, I., & Brown, M. E. 2015, AJ, 150, 174
Wong, I., & Brown, M. E. 2016, AJ, 152, 90
Yoshida, F. & Nakamura, T. 2005, AJ, 130, 2900
Yoshida, F. & Nakamura, T. 2007, Planet. Space Sci., 55, 1113
Yoshida, F. & Nakamura, T. 2008, PASJ, 60, 297
Yoshida, F., Nakamura, T., Watanabe, J.-I., et al. 2003, PASJ, 55, 701
Yoshida, F. & Terai, T. 2017, AJ, 154, 71
Yoshida, F., Terai, T., Ito, T., et al. 2019, Planet. Space Sci., 169, 78
Yoshida, F., Terai, T., Ito, T., et al. 2020, Planet. Space Sci., 190, 104977