A REPRESENTATION FORMULA FOR INDEFINITE IMPROPER 
AFFINE SPHERES

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Abstract. We construct a new representation formula for indefinite improper 
affine spheres in terms of two para-holomorphic functions and study singular-
ities which appear in this representation formula. As a result, it follows that 
cuspidal cross caps never appear as the singularities on indefinite improper 
affine spheres and so on. Comparison with other representation formulae are 
also studied.

Introduction

Affine spheres are important objects in the study of affine differential geometry. 
They have close relations to the theory of minimal surfaces, Monge-Ampère equa-
tions, Liouville equation, Tzitzeica equation and so on. In this paper, we construct 
a representation formula for indefinite improper affine spheres in the affine 3-space 
(Theorem 2.3), in terms of two para-holomorphic functions. This may be regarded 
as an indefinite version of the representation formula for locally strongly convex im-
proper affine spheres obtained by A. Martínez [15]. Here, indefinite improper affine 
spheres mean improper affine spheres with indefinite affine metric. It is found that 
affine spheres which are represented by these formulae may have singularities and 
we investigate them. As a result, we have peculiar examples of indefinite improper 
affine spheres with singularities (Section 5).

So far, various kinds of representation formulae for (improper) affine spheres 
are studied by several authors. Perhaps the earliest version, which is discovered 
in early twentieth century, is due to Blaschke [1]. In this representation formula, 
improper affine spheres are represented in terms of two smooth functions. Another 
representation formula is due to Cortés [2]. His formula represents special class of 
improper affine spheres which are related to special Kähler structure, in terms of 
one holomorphic function. The feature of his representation formula is that it even 
covers higher (even-) dimensional improper affine spheres.

Recently, A. Martínez made a representation formula for locally strongly convex 
improper affine spheres in terms of two holomorphic functions [15]. The feature 
of Martínez’ representation formula is that it includes even improper affine spheres 
with singularities. In the same paper, he also introduced the notion of IA-map, a 
class of (locally strongly convex) improper affine spheres with singularities which 
have close relations with his representation formula. He also studied a correspon-
dence between improper affine spheres and flat fronts in hyperbolic 3-space (see [6], 
[7]) in [14].

Following Martínez’ manner, we introduce generalized IA-maps in Section 2 
(Definitions 2.4, 2.5). The singularities on indefinite generalized IA-maps have 
different properties from those on locally strongly convex ones. The singularities 
on locally strongly convex generalized IA-maps are either singularity of fronts, or

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branch points. Contrarily, the singularities on indefinite generalized IA-maps, however, are not necessarily fronts nor branch points. Moreover, the cuspidal cross caps never appear as singularities of indefinite generalized IA-maps (Theorem 3.7), while they appear often in the case of spacelike maximal surfaces in Lorentzian 3-space. Furthermore, we found a strange example of singularity on indefinite generalized IA-map (Section 5). This is not $A$-equivalent to any singularities on improper affine spheres obtained by the projection of a generalized geometric solution of certain Monge-Ampère system which are considered in [10] (Section 3).

1. Preliminaries

Before explaining the main part of results, we would like to review some definitions and basic known facts about affine differential geometry, para-complex number and singularity theory.

1.1. Improper affine spheres. At first, we would like to introduce briefly the affine differential geometry. For the detailed exposition, see [13] and [18].

Let $(f, \xi)$ be a pair of an immersion $f : M^n \to R^{n+1}$ into the affine space $R^{n+1}$ and the vector field $\xi$ on $M^n$ along $f$ which is transversal to $f_*(TM)$. Then, the Gauss-Weingarten formula become

\[
\begin{align*}
D_X f_*(Y) &= f_*(\nabla_X Y) + g(X, Y)\xi, \\
D_X \xi &= -f_*(S X) + \tau(X)\xi,
\end{align*}
\]

where $D$ is the standard connection on $R^{n+1}$. Here, $g$ is called the affine metric of a pair $(f, \xi)$. It can be easily shown that the rank of the affine metric $g$ is invariant under the change of transversal vector field $\xi$. So we call $f$ a locally strongly convex immersion (respectively indefinite immersion) if $g$ is positive definite (resp. indefinite). Given an immersion $f : M^n \to R^{n+1}$, we can choose uniquely the transversal vector field $\xi$ which satisfy the following conditions,

1. $\tau \equiv 0$, (or equivalently $D_X \xi \in f_*(TM)$ for all $X \in \mathfrak{X}(M)$),
2. $\text{vol}_g(X_1, \ldots, X_n) = \det(f_*(X_1, \ldots, X_n, \xi))$ for all $X_1, \ldots, X_n \in \mathfrak{X}(M)$,

where $\text{vol}_g$ is the volume form of the (pseudo-)Riemannian metric $g$ and $\det$ is the standard volume element of $R^{n+1}$. The transversal vector field $\xi$ which satisfies above two conditions is called a Blaschke normal (or affine normal) and a pair $(f, \xi)$ of an immersion and its Blaschke normal is called a Blaschke immersion.

A Blaschke immersion $(f, \xi)$ with $S = 0$ in [18] is called an improper affine sphere. In this case $\xi$ becomes constant vector because $\tau = 0$. Hence hereafter, we can think of a transversal vector field $\xi$ of a improper affine sphere as $\xi = (0, \ldots, 0, 1)$ after certain affine transformation of $R^{n+1}$.

The conormal map $\nu : M^n \to (R^{n+1})^*$ for a given Blaschke immersion $(f, \xi)$ is defined as the immersion which satisfy the following conditions,

1. $\nu(f_*(X)) = 0$ for all $X \in \mathfrak{X}(M)$,
2. $\nu(\xi) = 1$.

For an improper affine sphere with the Blaschke normal $(0, \ldots, 1)$, we can write $\nu = (n, 1)$ with a smooth map $n : M^2 \to R^n$.

Using the notations defined as above, we can now state the representation formula for locally strongly convex improper affine spheres (possibly with singularities) by Martínez in [15]. In that paper, he define the notion of improper affine maps, a generalization of improper affine spheres which possibly have singularities (Definition 2.4) and give its representation formula as follows:
Theorem 1.1 (Theorem 3 in [15]). Let $\psi = (x, \varphi) : \Sigma^2 \to \mathbb{R}^2 \times \mathbb{R}$ be an improper affine map. Then there exists a regular complex curve $\alpha := (F, G) : \Sigma^2 \to \mathbb{C}^2$ such that,

$$\psi = \left( G + \bar{F}, \frac{1}{2}(|G|^2 - |F|^2) + \text{Re} \left( GF - 2 \int FdG \right) \right).$$

Here, the conormal map of $\psi$ becomes

$$\nu = (\bar{F} - G, 1).$$

Conversely, given a Riemann surface $\Sigma$ and a complex curve $\alpha := (F, G) : \Sigma \to \mathbb{C}^2$, then [12] gives an improper affine map which is well defined if and only if $\int FdG$ does not have real periods.

1.2. Para-complex numbers. The set of para-complex numbers $\tilde{\mathbb{C}}$, is an algebra over $\mathbb{R}$ which is defined as

$$\tilde{\mathbb{C}} := \{ a + jb | a, b \in \mathbb{R}, j^2 = 1 \}.$$

For $z = u + jv \in \tilde{\mathbb{C}}$ the conjugate $\bar{z}$ is defined as $\bar{z} := u - jv$ and absolute value $|z|$ is defined as $|z| := z\bar{z} = u^2 - v^2$. As an analogy to the holomorphic function, we can define the so-called para-holomorphic function: A map $F : \tilde{\mathbb{C}} \to \tilde{\mathbb{C}}$ is called a para-holomorphic function if $F$ satisfies the para-Cauchy-Riemann equation, that is,

$$\begin{cases}
\frac{\partial f^1}{\partial u} = \frac{\partial f^2}{\partial v}, \\
\frac{\partial f^1}{\partial v} = \frac{\partial f^2}{\partial u},
\end{cases}$$

where $F(u + jv) = f^1(u, v) + jf^2(u, v)$. The set of para-holomorphic functions forms an algebra. For a para-holomorphic function $F : \tilde{\mathbb{C}} \to \tilde{\mathbb{C}}$, we define its differential $F'$ as

$$F'(u + iv) := \frac{\partial f^1}{\partial u} + j\frac{\partial f^2}{\partial u},$$

(or equivalently, $F'(u + iv) := \frac{\partial f^2}{\partial v} + j\frac{\partial f^1}{\partial v}$)

where $F(u + jv) = f^1(u, v) + jf^2(u, v)$.

Since [13] is reduced to the wave equation

$$\frac{\partial^2 f_j}{\partial u^2} - \frac{\partial^2 f_j}{\partial v^2} = 0 \quad (j = 1, 2)$$

and the general solution of the wave equation, which is known as d’Alembert solution, is given by arbitrary two smooth functions, a para-holomorphic function is described in terms of two smooth functions. Concretely, for a para-holomorphic function $F$, there exist two smooth functions $\rho$ and $\sigma$ on $\mathbb{R}^2$, such that

$$F(u + jv) = \rho(u + v) + \sigma(u - v) + j \{ \rho(u + v) - \sigma(u - v) \}$$

holds. For more detailed exposition on para-complex numbers, see [8].

1.3. Criteria for singularities. Here, we explain criteria for singularities of smooth maps (see [3], [9] and [17]). The explanation is restricted to the case for smooth maps from 2-dimensional manifolds $M^2$ to the affine 3-space $\mathbb{R}^3$ because all the affine spheres are considered to be 2-dimensional in this paper.

Consider a smooth map $f : U \to \mathbb{R}^3$ from an open subset $U$ of $\mathbb{R}^2$ to the affine 3-space $\mathbb{R}^3$. A point $p$ on $U$ is called a singular point of $f$, if $f$ is not immersive at $p \in U$. On studying the local properties of singularity on smooth maps, one usually consider its map-germ. Let $f : U \to \mathbb{R}^3$ and $g : V \to \mathbb{R}^3$ are smooth maps from open sets $U, V$ of $\mathbb{R}^2$ to $\mathbb{R}^3$ respectively, with $p \in U \cap V$. Then we call $f$ and $g$ defines the same map-germ at $p$ if there exists an open set $W \subset \mathbb{R}^2$ with $p \in W$ and
\( W \subset U \cap V \) such that \( f = g \) holds on \( W \). Next, we introduce \( A \)-equivalence, an equivalence relation between two singularities. Given two smooth maps \( f : U \to R^3 \) and \( g : \hat{U} \to R^3 \), and assume \( p \in U \) and \( \hat{p} \in \hat{U} \) are singular points. Consider two map-germ \((f,p)\) around \( p \in U \) and \((g,\hat{p})\) around \( \hat{p} \in \hat{U} \). Then \((f,p)\) and \((g,\hat{p})\) is called \( A \)-equivalent if there exists two diffeomorphism-germ \( \varphi : U \to \hat{U} \) and \( \psi : R^3 \to R^3 \) such that \( \psi \circ f = \hat{f} \circ \varphi \) as a map-germ around \( p \in U \) and \( \hat{p} = \varphi(p) \) hold.

In this paper, we mainly consider the special class of singularities, that is, fronts and fronts. A smooth map (possibly with singularities) \( f : U \to R^3 \) from open subset \( U \subset R^2 \) to the affine 3-space \( R^3 \) is called a frontal (or a frontal map) if there exists a unit normal vector field \( \nu \) to \( f \) (even on singular points). This is equivalent to the existence of a map to the 2-sphere \( \nu : U \to S^2 \subset R^3 \) which satisfies
\[
\langle df, \nu \rangle = 0, \tag{1.4}
\]
where \( \langle , \rangle \) is the Euclidean metric of \( R^3 \). By the identification of the unit tangent bundle \( T_1 R^3 \) and the unit cotangent bundle \( T^*_1 R^3 \), frontals can be considered as the projection of a Legendrian map into \( T^*_1 R^3 \) with respect to the canonical contact structure as follows. Let \( f : U \to R^3 \) be a front and \( \nu \) be its unit normal vector field. Then \( f \) and \( \nu \) defines a map into the unit tangent bundle, \((f,\nu) : U \to T_1 R^3 \). We can consider this map as \((f,\nu) : U \to T^*_1 R^3 \) by the above identification.

Then the condition \((1.4)\) is equivalent to that \((f,\nu) \) is Legendrian map into \( T^*_1 R^3 \). Conversely, from a Legendrian map into \( T^*_1 R^3 \), we can make a frontal into \( R^3 \) by the projections into \( R^3 \) and \( S^2 \) respectively. A frontal \( f : U \to R^3 \) is called a front if \((f,\nu) \) defined as above is an immersion. From above arguments, it can be seen that the notion of fronts and frontals do not depend on the choice of the Riemannian metric \( g \) of \( R^3 \).

Let \( \det \) be the standard volume element of \( R^3 \). Then the function \( \lambda(u,v) := \det(f_u,f_v,\nu) \) is called the signed area density function where \( \nu u \) is the unit normal of \( f \). It is obvious from the definition that \( p \in U \) is a singular point if and only if \( \lambda(p) = 0 \) holds. A singular point \( p \in M^2 \) is called non-degenerate if the derivative \( d\lambda \) of the signed area density function does not vanish at \( p \).

The typical examples of fronts are cuspidal edges and swallowtails, while those of frontals are a cuspidal cross caps. Here, cuspidal edges, swallowtails and cuspidal cross caps are defined as a smooth map \( f : M^2 \to N^3 \) which are \( A \)-equivalent to the following maps, \( f_C \), \( f_S \) and \( f_{CCR} \) respectively:
\[
\begin{align*}
  f_C(u,v) & := (u^2, u^3, v), \quad f_S(u,v) := (3u^4 + u^2v, 4u^3 + 2uv, v), \\
  f_{CCR}(u,v) & := (u, v^2, uv^3).
\end{align*}
\]

In \cite{4} and \cite{9}, a geometric criteria for singular points to be cuspidal edges, swallowtails and cuspidal cross caps are given. To make the criterion, they define singular curve, null direction and so on. If \( p \in U \) is a non-degenerate singular point, then the implicit function theorem implies that the singular set \( \Sigma_f \) around \( p \) becomes locally a regular curve because \( \Sigma_f \) coincides with the zero sets of signed area density function \( \lambda \). This curve is called a singular curve and denoted by \( \gamma : (-\varepsilon,\varepsilon) \to U \). Usually, we fix the parametrization of \( \gamma \) as to \( \gamma(0) = p \). The singular direction of \( f \) at \( \gamma(t) \) is the 1-dimensional subspace of \( T_p U \) which is spanned by \( \gamma'(0) \in T_p U \). The null direction is defined as the kernel of the differential map \( f_* \). This is defined uniquely only in the case that image of the differential map \( f_* \) become 1-dimensional. A vector field \( \eta(t) \) along the singular curve \( \gamma(t) \) is called a null vector field if it associates the non-zero vector belonging to the null direction in \( T_{\gamma(t)} U \) to each \( t \).
Using the above definitions, criteria for cuspidal edges, swallow tails and cuspidal cross caps are given as follows.

Fact 1.2 (Criteria for cuspidal edges and swallowtails \[^4\]). Let \( p \) be a non-degenerate singular point of a front \( f \), \( \gamma \) the singular curve passing through \( p \), and \( \eta \) a null vector field along \( \gamma \). Then

1. \( p = \gamma(0) \) is a cuspidal edge if and only if the null direction and the singular direction are transversal, that is, \( \det(\gamma'(0), \eta(0)) \neq 0 \) holds, where \( \det \) denotes the determinant of \( 2 \times 2 \) matrices and where we identify the tangent space in \( T_pU \) with \( \mathbb{R}^2 \).
2. \( p = \gamma(0) \) is a swallowtail if and only if
   \[
   \det(\gamma'(0), \eta(0)) = 0 \quad \text{and} \quad \frac{d}{dt}\bigg|_{t=0} \det(\gamma'(t), \eta(t)) \neq 0
   \]
   hold.

Fact 1.3 (Criterion for cuspidal cross caps \[^4\]). Let \( f : U \to \mathbb{R}^3 \) be a frontal and \( \gamma(t) \) a singular curve on \( U \) passing through a non-degenerate singular point \( p = \gamma(0) \). We set
   \[
   \Psi(t) := \det(\tilde{\gamma}', D^f_\eta \nu, \nu),
   \]
   where \( \tilde{\gamma} = f \circ \gamma \), \( D^f_\eta \nu \) is the canonical covariant derivative along a map \( f \) induced from the standard connection on \( \mathbb{R}^3 \), and \( t' = dt/dt \). Then the germ of \( f \) at \( p = \gamma(0) \) is \( A \)-equivalent to a cuspidal cross cap if and only if

1. \( \eta(0) \) is transversal to \( \gamma'(0) \),
2. \( \Psi(0) = 0 \) and \( \Psi'(0) \neq 0 \).

2. Representation formula for indefinite improper affine spheres

In this section, we introduce a representation formula for indefinite improper affine spheres by modifying Martínez' method in \[^15\]. To do this, we should review a duality relation for improper affine spheres with indefinite affine metric. Throughout this paper, we fix its affine normal to \( \xi = (0, 0, 1) \) by the appropriate affine transformation of \( \mathbb{R}^3 \).

Proposition 2.1 (Duality relations). Let \( f : \Sigma^2 \to \mathbb{R}^3 \) be an indefinite improper affine sphere, \( \xi \) its affine normal and \( \nu : \Sigma^2 \to \mathbb{R}^3 \) the conormal map of \( f \), where \( \Sigma^2 \) is an open subset of \( \mathbb{R}^2 \). Take a coordinate system \((u, v)\) of \( \Sigma^2 \) so that the affine metric \( g \) of \( f \) (which is indefinite in this case) is represented as \( g = E(du^2 - dv^2) \). Then, the following identities hold:
   \[
   \begin{align*}
   f_u &= \nu \times \nu_v \\
   f_v &= \nu \times \nu_u \\
   \nu_u &= f_v \times \xi \\
   \nu_v &= f_u \times \xi
   \end{align*}
   \]
   where \( \times \) is the outer product on \( \mathbb{R}^3 \), \( f_u := \frac{\partial f}{\partial u} \) and \( f_v := \frac{\partial f}{\partial v} \).

Proof. The duality relations for affine spheres are well-known for the other kinds of affine spheres. The proof is almost parallel to them. In fact, we can obtain the duality relations for indefinite improper affine spheres by taking care of the definition of conormal maps and the Gauss-Codazzi equations. The case of definite improper affine spheres is presented in \[^17\]. \( \square \)

Let \( \psi = (x, \varphi) : \Sigma^2 \to \mathbb{R}^3 (= \mathbb{R}^2 \times \mathbb{R}) \) be an indefinite improper affine sphere with the Blaschke normal \( \xi = (0, 0, 1) \), where \( x : \Sigma^2 \to \mathbb{R}^2 \) and \( \varphi : \Sigma^2 \to \mathbb{R} \) are smooth maps. Then the conormal vector field \( \nu : \Sigma^2 \to \mathbb{R}^3 (= \mathbb{R}^2 \times \mathbb{R}) \) can be represented as \( \nu = (n, 1) \) and the affine metric is written as \( g = -\langle dx, du \rangle \), where \( \langle , \rangle \) is the standard Euclidean inner product of \( \mathbb{R}^2 \).
Using Proposition 2.1, we obtain the following characterization of indefinite improper affine spheres.

**Proposition 2.2.** Define \( \tilde{L}_\psi : \Sigma^2 \to \tilde{C}^2 \) as
\[
\tilde{L}_f := x + jn,
\]
where \( x \) and \( n \) are defined as above, that is, \( x := \pi \circ \psi \) and \( n := \pi \circ \nu \) with the projection \( \pi : R^3 \to R^2 \). Identify \( \tilde{C}^2 \) with \( R^4 \) as,
\[
\tilde{C}^2 \ni z_1 = y_0 + jy_1, z_2 = y_2 + jy_3 \longleftrightarrow (y_0, y_1, y_2, y_3) \in R^4
\]
Then \( \tilde{L}_\psi \) annihilates 2-forms \( dy_0 \wedge dy_1 + dy_2 \wedge dy_3 \) and \( dy_0 \wedge dy_2 + dy_1 \wedge dy_3 \), in other words,
\[
\begin{align*}
\left\{ \begin{array}{l}
dy_0 \wedge dy_1 + dy_2 \wedge dy_3 |_{\tilde{L}_\psi(\Sigma^2)} = 0, \\
dy_0 \wedge dy_2 + dy_1 \wedge dy_3 |_{\tilde{L}_\psi(\Sigma^2)} = 0.
\end{array} \right.
\end{align*}
\]
(2.1) hold.

**Proof.** Applying Proposition 2.2, we have
\[
\tilde{L}_\psi^*(dy_0 \wedge dy_1 + dy_2 \wedge dy_3) = dx^1 \wedge dx^2 + dn^1 \wedge dn^2
\]
\[
= \left\{ (x^1_n x^2_n - x^2_n x^1_n) + (n^1_n n^2_n - n^1_n n^2_n) \right\} du \wedge dv
\]
\[
= 0
\]
and
\[
\tilde{L}_\psi^*(dy_0 \wedge dy_2 + dy_1 \wedge dy_3) = dx^1 \wedge dn^1 + dx^2 \wedge dn^2
\]
\[
= \left\{ (x^1_n n^1_n - x^2_n n^1_n) + (x^2_n n^2_n - x^2_n n^2_n) \right\} du \wedge dv
\]
\[
= 0
\]
where \( x = (x^1, x^2) \) and \( n = (n^1, n^2) \). \( \square \)

Proposition 2.2 immediately leads us to the following representation formula.

**Theorem 2.3.** (1) Let \( \psi = (x, \varphi) : \Sigma^2 \to R^2 \times R \) be an indefinite improper affine sphere whose affine normal is fixed to \( \xi = (0, 0, 1) \). Let \( \nu = (n, 1) : \Sigma^2 \to R^2 \times R \) be the conormal map of \( \psi \). Then there exists a para-holomorphic curve \( \alpha := (F, G) : \Sigma^2 \to \tilde{C}^2 \) such that \( |dF| \neq |dG| \) and
\[
\left\{ \begin{array}{l}
x = F - \bar{G}, \\
n = F + G.
\end{array} \right.
\]
(2.2) hold everywhere.

(2) Conversely, given a para-holomorphic curve \( (F, G) : \Sigma \to \tilde{C}^2 \), then
\[
\psi = \left( x, - \int \langle n, dx \rangle \right)
\]
gives an indefinite improper affine sphere (possibly with singularities), where \( x \) and \( n \) are defined by 2.2. Moreover, if we write \( F \) and \( G \) as \( F = \)
\( f + jG, G = g + jg^2 \), then \( \psi \) can be expressed as

\[
\psi = \left( f^1 - g^1, f^2 + g^2, \right.
\]

\[
- \int \left\{ (f^1 + g^1)(f^1_u - g^1_u) + (-f^2 + g^2)(f^2_u + g^2_u) \right\} \, du - \int \left\{ (f^1 + g^1)(f^2_u - g^2_u) + (-f^2 + g^2)(f^1_u + g^1_u) \right\} \, dv
\]

**Proof.**

(1) First, define \( F : \Sigma \to \tilde{C} \) as

\[
F := f^1 + jf^2 := \frac{x^1 + y^1}{2} + j\frac{x^2 - y^2}{2}. 
\]

Then, by Proposition 2.2 we have

\[
\begin{align*}
  f^1_u &= f^2_v, \\
  f^2_u &= f^1_v
\end{align*}
\]

and \( F \) is a para-holomorphic function on \( \Sigma^2 \) with respect to the para-complex structure induced from \( g \).

Take a new coordinate system \( \{w_1, w_2\} \) on \( \tilde{C}^2 \) as,

\[
\begin{align*}
  w_1 &= y_0 + y_1 + j\frac{y_2 - y_3}{2}, \\
  w_2 &= -\frac{y_0 + y_1}{2} + j\frac{y_2 + y_3}{2}.
\end{align*}
\]

Then, (2.4) is equivalent to

\[
dw_1 \wedge dw_2 |_{I_{\tilde{L}}(\Sigma)} = 0.
\]

Here \( w_1 = F \) and \( w_1 \) induce the same para-complex structure on \( \Sigma \) as the one induced from \( g \). Thus (2.4) implies that \( w_2 \) also defines a para-holomorphic function on \( \Sigma \). Define a para-holomorphic function \( G : \Sigma \to \tilde{C} \) as \( G := w_2 \), then we have (2.2). The condition that \( |dF| \neq |dG| \) is equivalent to the absence of singularity on \( \psi \).

(2) We can check that \( \psi \) defines an indefinite improper affine sphere by direct calculation. However, it is almost obvious that \( \psi \) defines an indefinite improper affine sphere by the following reason. For \( \psi : \Sigma^2 \to R^3 \), a duality relation is equivalent to the condition to be an affine sphere. On the other hand, a duality relation is equivalent to para-Cauchy-Riemann equation in this case.

Next, we introduce the notion of IA-maps and generalized IA-maps, which is a generalization of improper affine sphere. IA-maps are firstly defined by A. Martínez in [15], as a generalization of improper affine sphere which permits a singularity. Modifying the Martinez’ definition of IA-maps, we define generalized IA-maps as the following.

**Definition 2.4** (Definition 1 in [15]). A map \( \psi = (x, \varphi) : \Sigma^2 \to R^3 (= R^2 \times R) \), where \( x : \Sigma^2 \to R^2 \) and \( \varphi : \Sigma^2 \to R \) are smooth maps, is called an IA-map (respectively a generalized IA-map) if \( \tilde{L}_\psi : \Sigma^2 \to C^2 \), which is defined by

\[
\tilde{L}_\psi := x + \sqrt{-1} n,
\]

becomes a SL-immersion (resp. SL-map) with respect to the symplectic structure \( \omega' \) and the calibration \( \text{Re}(\sqrt{-1} \Omega') \), where \( \omega' = \frac{\sqrt{-1}}{2} (d\zeta_1 \wedge d\bar{\zeta}_1 + d\zeta_2 \wedge d\bar{\zeta}_2) \) and \( \Omega' = d\zeta_1 \wedge d\zeta_2 \).
As an analogy to this definition, it is appropriate to define the class of indefinite improper affine spheres with singularities as follows.

**Definition 2.5.** A map $\psi = (x, \varphi) : \Sigma^2 \to \mathbb{R}^3(= \mathbb{R}^2 \times \mathbb{R})$, where $x : \Sigma^2 \to \mathbb{R}^2$ and $\varphi : \Sigma^2 \to \mathbb{R}$ are smooth maps, is called an *indefinite IA-map* (respectively, an *indefinite generalized IA-map*) if $\tilde{\psi}$ and $\varphi$ are locally strongly convex IA-maps.

From now on, to avoid the confusion, we call an IA-map (respectively, a generalized IA-map) with positive definite affine metric an *locally strongly convex IA-maps* (resp. an *locally strongly convex generalized IA-map*).

### 3. Singularities of IA-maps

The representation formula for locally strongly convex improper affine spheres and that for indefinite ones look very similar. However, singular points which appear in two representation formulae have quite different properties each other. Concretely, in the locally strongly convex case, except for on branch points, all the singular points are fronts. On the other hand, in the indefinite case, a singular point is not necessarily a front even if it is not a branch point. Indeed, we can easily find examples of frontal maps which are not fronts in indefinite case by direct calculations.

#### 3.1. Locally strongly convex case.

To study the properties of singularities on a locally strongly convex IA-map $\psi : \Sigma^2 \to \mathbb{R}^3$, we will first check the condition for $p \in \Sigma^2$ to be a singular point of $\psi$, which is also given in [15].

**Proposition 3.1 (15).** Let $\psi : \Sigma \to \mathbb{R}^3$ be an locally strongly convex generalized IA-map. From the representation formula (1.2) for locally strongly convex IA-maps in Theorem 1.1, $p \in \Sigma^2$ is a singular point of $\psi$ if and only if $|dF| = |dG|$ holds on $p \in \Sigma^2$.

**Proof.** From the explicit form of the representation formula (2.3), the differentiation of $\psi$ becomes

$$
\psi_u = (f_u^1 + g_u^1)(0,1,g^1) - f_u^1(g_u^1 + f_u^1) + (g_u^2 + f_u^2)(g_u^2 - f_u^2),
$$

$$
\psi_v = (-f_u^1 - g_u^1)(-g^1 + f^1)(g_u^2 + f_u^2) + (g_u^2 + f_u^2)(g_u^1 - f_u^1).
$$

This can be rewritten as

$$
\psi_u = (f_u^1 + g_u^1)(1,0,g^1 - f^1) + (g_u^2 - f_u^2)(0,1,g^2 + f^2)
$$

$$
\psi_v = -(g_u^2 + f_u^2)(1,0,g^1 - f^1) + (g_u^1 - f_u^1)(0,1,g^2 + f^2).
$$

The linearly independence of the above two vectors implies that the condition to be a singular points is

$$
(f_u^1 + g_u^1)(g_u^1 - f_u^1) - (g_u^2 - f_u^2)(-g_u^2 + f_u^2) = 0.
$$

This is equivalent to $|dF| = |dG|$. \hfill \Box

Next we show that the singularity of locally strongly convex improper affine spheres is always locally a front unless it is a branch point, that is, $\psi_u = \psi_v = 0$ holds on that point. In other words, Legendrian lift $L_\psi = (\psi, \tilde{\nu}) : \Sigma^2 \to T\mathbb{R}^3$ is immersive on $\Sigma^2$ (even on the singular point of $\psi$), where $\tilde{\nu}$ is the unit normal of $\psi$. 

Proposition 3.2. Singularity of a locally strongly convex generalized IA-map is always locally a front if \(|dF| = |dG| = 0\) does not hold.

Proof. The proof follows from the direct calculation. We check that \(\text{Ker}\psi \neq \text{Ker}\tilde{\nu}\) holds at singular points of \(\psi\).

From Theorem [3.1], the explicit representation of \(\tilde{\nu}\) is

\[
\tilde{\nu} = \frac{1}{(1 + (f^1 - g^1)^2 + (f^2 + g^2)^2)^2} (f^1 - g^1, -(f^2 + g^2), 1),
\]

and its differentiations are

\[
\tilde{\nu}_u = \frac{1}{\sqrt{1 + (f^1 - g^1)^2 + (f^2 + g^2)^2}} \left( (f^1_u - g^1_u)(1 + (f^2 + g^2)^2) - (f^2 + g^2)(f^1 - g^1)(f^2_u + g^2_u),
\right.
\]

\[
- (f^2_u + g^2_u)(1 + (f^1 - g^1)^2) + (f^2 + g^2)(f^1 - g^1)(f^1_u - g^1_u),
\]

\[
- (f^1 - g^1)(f^1_u - g^1_u) - (f^2 + g^2)(f^2_u + g^2_u) \right),
\]

\[
\tilde{\nu}_v = \frac{1}{\sqrt{1 + (f^1 - g^1)^2 + (f^2 + g^2)^2}} \left( - (f^2_u - g^2_u)(1 + (f^2 + g^2)^2) - (f^2 + g^2)(f^1 - g^1)(f^2_u + g^2_u),
\right.
\]

\[
- (f^1_u + g^1_u)(1 + (f^1 - g^1)^2) - (f^2 + g^2)(f^1 - g^1)(f^2_u - g^2_u),
\]

\[
(f^1 - g^1)(f^2_u - g^2_u) - (f^2 + g^2)(f^1_u + g^1_u) \right).
\]

The proof of Proposition [3.1] implies that

\[
(g^2_u + f^2_u)\tilde{\nu}_u + (g^1_u + f^1_u)\tilde{\nu}_v = 0
\]

holds at a singular point of \(f\). Then, to prove \(L_\psi\) is an immersion, we have only to show that

\[
(g^2_u + f^2_u)\tilde{\nu}_u + (g^1_u + f^1_u)\tilde{\nu}_v \neq 0
\]

holds at any singular point.

Here,

\[
\sqrt{1 + (f^1 - g^1)^2 + (f^2 - g^2)^2}\left\{ (g^2_u + f^2_u)\tilde{\nu}_u + (g^1_u + f^1_u)\tilde{\nu}_v \right\}
\]

\[
= (- (g^1 - f^1)(g^2 + f^2)((g^1_u + f^1_u)^2 + (g^2_u + f^2_u)^2) + 2(1 + (f^1 - g^1)^2)(f^1_u g^2_u - f^2_u g^1_u),
\]

\[
- (1 + (f^1 - g^1)^2)(g^1_u + f^1_u)^2 + (g^2_u + f^2_u)^2) + 2(f^1 - g^1)(f^2 + g^2)(f^1 u g^2_u - f^2 u g^1_u),
\]

\[
- (f^2 + g^2)((f^1_u + g^1_u)^2 + (g^2_u + f^2_u)^2) - 2(f^1 - g^1)(f^1 u g^2_u - f^2 u g^1_u))
\]

\[
= (- (g^1 - f^1)(g^2 + f^2)A + 2(1 + (f^1 - g^1)^2)B, -(1 + (f^1 - g^1)^2)A + 2(f^1 - g^1)(f^2 + g^2)B,
\]

\[
- (f^2 + g^2)A - 2(f^1 - g^1)B).
\]

where, \(A = (f^1_u + g^1_u)^2 + (g^2_u + f^2_u)^2\) and \(B = (f^1 u g^2_u - f^2 u g^1_u)\). Here \(A > 0\) because \(|dF| = |dG| = 0\) does not hold. If

\[
\sqrt{1 + (f^1 - g^1)^2 + (f^2 - g^2)^2}\left\{ (g^2_u + f^2_u)\tilde{\nu}_u + (g^1_u + f^1_u)\tilde{\nu}_v \right\} = 0
\]

holds, then the computation of \(B/A\) by using the fact that each element becomes 0 gives \(1 + (f^1 - g^1)^2 + (f^2 + g^2)^2 = 0\). However, this is contradiction. So

\[
(g^2_u + f^2_u)\nu_u + (g^1_u + f^1_u)\nu_v \neq 0
\]
holds at any point (even at a singular point)
This completes the proof of Proposition 3.2. \( \square \)

Remark 3.3. We can state the relation between \( \psi \), \( \tilde{L}_\psi \) and \( L_\psi \) as follows. \( L_\psi \)
is an immersion if and only if \( \tilde{L}_\psi \) is an immersion. So a locally strongly convex
generalized IA-map \( \psi \) is a front if and only if \( \tilde{L}_\psi \) is an immersion.
On the other hand, if \( \tilde{L}_\psi \) is not immersive, then \( \psi_u = \psi_v = 0 \) holds at that point.
On such a point \( |dF| = |dG| = 0 \) holds. We call such a point a branch point. Note
that branch points are isolated because \( F \) and \( G \) are both holomorphic functions.
Above fact can be considered as another proof of Proposition 3.2.

### 3.2. Indefinite case

So far in this section, we consider the singularity of locally strongly convex IA-maps. Unlike the locally strongly convex case, any indefinite improper affine sphere (with singularities) is not necessarily a front. Indeed, we can construct several concrete examples which are frontal maps, but not fronts. So we shall check the property of singularity of indefinite improper affine spheres in the rest of this section.

First, similar to the locally strongly convex case, we check the condition for a pair of para-holomorphic functions so that the corresponding indefinite improper affine sphere has a singular point.

**Proposition 3.4.** Let \( \psi : \Sigma^2 \to \mathbb{R}^3 \) be an indefinite generalized IA-map Then, \( p \in \Sigma^2 \) is a singular point of \( \psi \) if and only if \( |dF| = |dG| \) holds on \( p \in \Sigma^2 \).

**Proof.** The proof is almost the same as that of Theorem 3.1 except for the detailed calculation.

From the explicit form of the representation formula \( \Sigma^2 \), the differentiation of \( \psi \) becomes
\[
\psi_u = (f_1^1 - g_1^1, f_1^2 + g_2^2, -(f_1^1 + g_1^1)(f_1^2 - g_2^2) - (-f_2^1 + g_2^1)(f_1^1 - g_1^2)),
\]
\[
\psi_v = (f_2^1 - g_2^1, f_1^1 + g_1^1, -(f_2^1 + g_1^2)(f_1^2 - g_2^2) - (-f_2^1 + g_2^1)(f_1^1 - g_1^1)).
\]

This can be rewritten as
\[
\psi_u = (f_1^1 - g_1^1)(1, 0, -(f_1^1 + g_1^1)) + (f_1^2 + g_2^2)(0, 1, f_2^1 - g_2^1)
\]
\[
\psi_v = (f_2^1 - g_2^1)(1, 0, -(f_2^1 + g_1^2)) + (f_1^1 + g_1^1)(0, 1, f_2^1 - g_2^1).
\]

The linearity independence of the above two vectors implies that the condition to be a singular points is
\[
(f_1^1 - g_1^1)(f_1^2 + g_2^2) - (f_2^1 + g_1^2)(f_1^2 - g_2^2) = 0.
\]

This is equivalent to \( |dF| = |dG| \).

As mentioned above, there exist frontal maps which are not fronts in the representation formula \( \Sigma^2 \). Next, we shall check the condition for an indefinite improper affine sphere to have this type of singularity.

**Proposition 3.5.** For an indefinite generalized IA-map \( \psi : \Sigma^2 \to \mathbb{R}^3 \), the following two conditions are equivalent:

1. \( \psi \) is a frontal map which is not a front at \( p \in \Sigma^2 \) and \( p \) is not a branch point.
2. One of the following two conditions are satisfied at \( p \in \Sigma^2 \).
   a. \( f_1^1 = f_2^2 \) and \( g_1^1 = g_2^2 \).
   b. \( f_1^1 = -f_2^2 \) and \( g_1^1 = -g_2^2 \).

**Proof.** From \( \Sigma^2 \), we already know that at any point, the tangent space of any indefinite generalized IA-map is spanned by two vectors \((1, 0, -(f_1^1 + g_1^1))\) and \((0, 1, f_2^1 - g_2^1)\). This means every indefinite generalized IA-map has an unit normal,
which is parallel to \((f^1 + g^1, -f^2 + g^2, 1)\). So every indefinite generalized IA-map is a frontal map. Therefore the condition (1) is satisfied if and only if \(\psi\) is not a front, i.e., \(L_{\psi}\) is not a immersion. Thus we have only to check the condition for \(L_{\psi}\) not to be an immersion.

From Theorem 2.3, the unit normal \(\tilde{\nu}\) of \(\psi\) is
\[
\tilde{\nu} = \frac{1}{\sqrt{\Delta}}(f^1 + g^1, -f^2 + g^2, 1)
\]
where \(\Delta = (f^1)^2 + (f^2)^2 + (g^1)^2 + (g^2)^2 + 2f^1g^1 - 2f^2g^2 + 1\).
Then, we have (3.5)
\[
\tilde{\nu}_u = -\frac{\Delta_u}{2\sqrt{\Delta}}(f^1 + g^1, -f^2 + g^2, 1) + \frac{1}{\sqrt{\Delta}}(f^1_u + g^1_u, -f^2_u + g^2_u, 0)
\]
\[
= \frac{1}{2\sqrt{\Delta}}\left\{ -\Delta_u(f^1 + g^1, -f^2 + g^2, 1) + 2\Delta(f^1_u + g^1_u, -f^2_u + g^2_u, 0) \right\}
\]
\[
= \frac{1}{2\sqrt{\Delta}}\left\{ -2(f^2_u - g^2_u)(f^2 - g^2) + (f^1 + g^1)^2 + 2(-f^2_u + g^2_u)(f^2 - g^2) + 2(f^1_u + g^1_u),
\right.
\]
\[
\left. -2(f^1 + g^1) + 2(f^2_u - g^2_u)(f^2 - g^2) \right\},
\]
\[
\tilde{\nu}_v = -\frac{\Delta_v}{2\sqrt{\Delta}}(f^1 + g^1, -f^2 + g^2, 1) + \frac{1}{\sqrt{\Delta}}(f^2_u + g^2_u, -f^1_u + g^1_u, 0)
\]
\[
= \frac{1}{2\sqrt{\Delta}}\left\{ -\Delta_v(f^1 + g^1, -f^2 + g^2, 1) + 2\Delta(f^2_u + g^2_u, -f^1_u + g^1_u, 0) \right\}
\]
\[
= \frac{1}{2\sqrt{\Delta}}\left\{ -2(f^1_u - g^1_u)(f^1 - g^1) + 2(f^2_u + g^2_u)(f^2 - g^2) + 2(f^2_u + g^2_u),
\right.
\]
\[
\left. -2(f^1_u + g^1_u)(f^1 - g^1) - 2(f^2_u - g^2_u)(f^2 - g^2) \right\},
\]
First, we consider the case that neither \(f^1_u + g^1_u, f^1_u - g^1_u, f^2_u + g^2_u\) nor \(f^2_u - g^2_u\) does not equal to 0. Since \(|dF| = |dG|\) holds on a singular point of \(\psi\), (3.3) implies that (3.6)
\[
\psi_u = \frac{f^1_u - g^1_u}{f^2_u - g^2_u} \tilde{\nu}_v = \frac{f^2_u + g^2_u}{f^1_u + g^1_u} \tilde{\nu}_v
\]
also holds on a singular point of \(\psi\). On the other hand, by the same reason, (3.6) implies that (3.7)
\[
\tilde{\nu}_v = \frac{f^2_u - g^2_u}{f^1_u - g^1_u} \tilde{\nu}_u = \frac{f^1_u + g^1_u}{f^2_u + g^2_u} \tilde{\nu}_v
\]
holds on a singular point of \(\psi\).

By comparing (3.6) and (3.7), and taking care of the fact that \(\psi\) is not a front if and only if \((L_{\psi})_u\) and \((L_{\psi})_v\) are parallel, we have condition (2) after elementary calculation.

In the case that either \(f^1_u + g^1_u, f^1_u - g^1_u, f^2_u + g^2_u\) or \(f^2_u - g^2_u\) equals 0, we have condition (2) by taking care of the fact that \(\psi\) is not a front at \(p\) and that \(p\) is not a branch point.

\[\square\]

Remark 3.6. In the same manner as Remark 3.3, we have the following relation between \(\psi\), \(L_{\psi}\) and \(L_{\phi}\) for an indefinite generalized IA-map.

\(L_{\psi}\) is an immersion if and only if \(L_{\psi}\) is an immersion. So an indefinite generalized IA-map \(\psi\) is a front if and only if \(L_{\psi}\) is an immersion.
On the other hand, if $L_\psi$ is not an immersion, then $|dF| = |dG| = 0$ hold and it follows that $\psi$ is a frontal which is not a front or $p$ is a branch point.

Now, we end this section by remarking that cuspidal cross caps never appear as the singularities of indefinite generalized IA-maps. This fact is proved by using the criterion for cuspidal cross caps (Fact 1.3) and the condition for indefinite generalized IA-map to be a frontal but not a front (Proposition 3.5).

\textbf{Theorem 3.7.} Let $\psi : \Sigma^2 \to \mathbb{R}^3$ be an indefinite generalized IA-map and $p \in \Sigma^2$ be a point. Then the germ of $\psi$ at $p$ is not $\mathcal{A}$-equivalent to the cuspidal cross cap.

\textit{Proof.} From Fact 1.3, it is sufficient to check that $\Psi'(0) = 0$. Here, we have only to consider the case of $f_1' = f_2'$ and $g_1' = g_2'$ because the proof of another case is parallel to it. Since $\gamma'(0) = (1, 1)$, $\eta(0) = (1, -1)$ and $\nu_{uu}(p) = \nu_{vv}(p)$, we have $\Psi'(0) = 0$. \hfill $\Box$

\textit{Remark 3.8.} Theorem 3.7 is proved in [11] for another setting. In [11], Ishikawa and Machida considered singularities on improper affine spheres given by projection derived from the framework of differential systems. They also have the same results as Theorem 3.7. In Section 5, we will discuss the relationship between the singularities on generalized indefinite IA-maps and those which appear in [11].

\section{4. Comparison with other representation formulae}

In addition to Martínez type representation formulae, there are several representation formulae for improper affine spheres. In this section, we compare the representation formulae for (mainly indefinite) improper affine spheres which are obtained in the previous sections, to the other representation formulae for improper affine spheres which are studied in [1], [2], [3] and [10]. The idea that relates the Martínez’ representation formula to the Cortés’ representation formula is originally due to T. Kurose ([11]).

At first, the relation between Martínez type representation formula and that obtained by Cortés-Lawn-Schäfer is as follows.

\textit{Remark 4.1.} For 2-dimensional case, we can derive Cortés-Lawn-Schäfer representation formula in [3] from Theorem 2.3 by taking $F = \frac{1}{2}(z-jf'(z)), G = \frac{1}{2}(z+jf'(z))$.

This fact is the analogy for the case of locally strongly convex ones as below (Remark 4.2).

The following two remarks are pointed out by T. Kurose.

\textit{Remark 4.2 ([11]).} For 2-dimensional case, we can derive Cortés’ representation formula in [2] from Martínez’ representation formula in [15] by taking $F = \frac{1}{2}(z-\sqrt{-1}f'(z)), G = \frac{1}{2}(z+\sqrt{-1}f'(z))$.

\textit{Remark 4.3.} By the identification of a para-holomorphic function to a pair of smooth functions explained in Section 1, we can derive Blaschke’s representation formula from the Martínez type representation formula in Theorem 2.3.

Here, two para-holomorphic functions $F,G$ can be written as

\[
\begin{align*}
F(u,v) &= \rho_1(u+v) + \sigma_1(u-v) + j\rho_2(u+v) - \sigma_1(u-v), \\
G(u,v) &= \rho_2(u+v) + \sigma_2(u-v) + j\rho_1(u+v) - \sigma_2(u-v)
\end{align*}
\]

by four smooth function $\rho_1, \rho_2, \sigma_1, \sigma_2$. Therefore if we take

\[
\begin{align*}
U_1 &= \rho_1 + \rho_2, \quad V_1 = \sigma_1 + \sigma_2, \\
U_2 &= -\rho_1 + \rho_2, \quad V_2 = \sigma_1 - \sigma_2,
\end{align*}
\]

then we have the Blaschke’s representation formula in [11].
5. Examples

In [15], Martínez gave several examples of locally strongly convex improper affine spheres by substituting holomorphic functions to his representation formula. Here, we construct indefinite improper affine spheres in the same manner as [15].

Example 5.1. Let \( (F, G) = (z^2, z^3) \), then Theorem 2.3 gives a concrete example of indefinite generalized IA-maps, which is a frontal map but not a front.

For this example, we have
\[
\begin{align*}
  f^1 &= u^2 + v^2, \\
  f^2 &= 2uv, \\
  g^1 &= u^3 + 3uv^2, \\
  g^2 &= 3u^2v + v^3,
\end{align*}
\]
and this implies
\[
\lambda = (f_u^1)^2 - (f_u^2)^2 - (g_u^1)^2 + (g_u^2)^2 = (u^2 - v^2)(4 - 9(u^2 - v^2))
\]
and
\[
\begin{align*}
  \lambda_u &= 4u(2 - 9u(u^2 - v^2)), \\
  \lambda_v &= -4v(2 - 9u(u^2 - v^2)).
\end{align*}
\]
Therefore, singular locus is \( \{(u, v) \in \mathbb{R}^2 | (u^2 - v^2)(4 - 9(u^2 - v^2)) = 0 \} \). Among them, only the origin is degenerate singular point, and \( \psi \) is a frontal but not a front on \( \{u^2 - v^2 = 0\} \) (they are all corank 1 map-germs except for on the origin.) while \( \psi \) is a frontal on \( C := \{4 - 9(u^2 - v^2) = 0\} \) (Figure 1).

By using the criteria for cuspidal edges and swallowtails, we can completely classify the singularities on \( C \). In fact, at any point on \( C \) except for the point \((-\frac{2}{3}, 0)\), singularities are \( \mathcal{A} \)-equivalent to the cuspidal edges while swallowtail appears at \((-\frac{2}{3}, 0)\). This assertion is verified as below.

Let \( C_1 := C \cap \{u \geq 0\} \) and \( C_2 := C \cap \{u \leq 0\} \). First, we show that the singularity is \( \mathcal{A} \)-equivalent to the cuspidal edge at any point on \( C_1 \). \( C_1 \) is parameterized as \( \gamma(t) := \frac{t}{3}(\cosh t, \sinh t) \). Since
\[
\begin{align*}
(f^2 + g^2)u(\gamma(t)) &= \frac{1}{3}\sinh t(1 + 2\cosh t), \\
(f^2 + g^2)v(\gamma(t)) &= \frac{2}{3}(\cosh t + 1)(2\cosh t - 1),
\end{align*}
\]
the null vector field at \( \gamma(t) \) is parallel to \((-\cosh t + 1)(2\cosh t - 1), \sinh t(1 + 2\cosh t)\)).

Therefore
\[
\det(\gamma'(t), \eta(t)) = -(\cosh t + 1)
\]
and it follows that the null vector field and the singular direction are transversal on \( C_1 \). \( C_2 \) is parameterized as \( \gamma(t) := -\frac{2}{3}(\cosh t, \sinh t) \) and
\[
\begin{align*}
(f^2 - g^2)u(\gamma(t)) &= -\frac{4}{3}(\cosh t + 1)(2\cosh t - 1), \\
(f^2 - g^2)v(\gamma(t)) &= -\frac{2}{3}\sinh t(1 + 2\cosh t),
\end{align*}
\]
hold. Thus the null vector field at \( \gamma(t) \) is parallel to \((-\sinh t(1 + 2\cosh t), (\cosh t + 1)(2\cosh t - 1)\)). Therefore
\[
\det(\gamma'(t), \eta(t)) = -\sinh t
\]
and it follows that the null vector field and the singular direction are transversal on \( C_2 \) except for on the point \((-\frac{2}{3}, 0)\) and that \( \det|_{t=0}(\gamma'(t), \eta(t)) \neq 0 \) on \((-\frac{2}{3}, 0)\).

Example 5.2. Let \( (F, G) = (z^3, z^4) \), then Theorem 2.3 also gives a concrete example of indefinite generalized IA-maps, which is a frontal map but not a front.

For this example, we have
\[
\begin{align*}
  f^1 &= u^3 + 3uv^2, \\
  f^2 &= 3u^2v + v^3, \\
  g^1 &= u^4 + 6u^2v^2 + v^4, \\
  g^2 &= 4u^3v + 4uv^3,
\end{align*}
\]
and this implies
\[\lambda = (f_u^1)^2 - (f_u^2)^2 - (g_u^1)^2 + (g_u^2)^2 = (u^2 - v^2)^2(9 - 16(u^2 - v^2))\]
and
\[
\begin{cases}
\lambda_u &= 12u(u^2 - v^2)(3 - 8(u^2 - v^2)), \\
\lambda_v &= -12v(u^2 - v^2)(3 - 8v(u^2 - v^2)).
\end{cases}
\]
Therefore, singular locus is \(\{(u, v) \in \mathbb{R}^2 | (u^2 - v^2)^2(9 - 16(u^2 - v^2)) = 0\}\). Among them, \(\{u^2 - v^2 = 0\}\) is the set of degenerate singular point. Moreover, \(\psi\) is a frontal but not front on \(\{u^2 - v^2 = 0\}\) (they are all corank 1 map-germs except for on the origin.) while \(\psi\) is a front on \(\tilde{C} := \{9 - 16(u^2 - v^2) = 0\}\). Above calculation shows that \((1,1)\) is the point where the following three conditions are satisfied, that is,

(i) \((1,1)\) is a degenerate singular point.
(ii) \(\psi\) is a frontal but not a front of corank 1 on \((1,1)\).
(iii) the set of degenerate singular point around \((1,1)\) is locally a smooth curve.

It is worth mentioning that any map-germs satisfying above conditions which appear in Ishikawa-Machida’s formulation is not \(A\)-equivalent to this example (Figure 2). We prove this fact in Section 6.

As similar to the Example 5.1, the singularities on \(\tilde{C}\) is completely classified by the same way. At any point on \(\tilde{C}\) except for the point \((-\frac{2}{3}, 0)\), singularity is \(A\)-equivalent to the cuspidal edge while swallowtail appears at \((-\frac{2}{3}, 0)\).

6. COMPARISON WITH FORMULATION BY ISHIKAWA-MACHIDA

As mentioned in Remark 3.8, Ishikawa and Machida also studied improper affine spheres with singularities in another setting. They considered (generalized) geometric solutions of a certain Monge-Ampère system \(\mathcal{M}\) on \(\mathbb{R}^5\) and studied the singularities of the projections of such (generalized) geometric solutions. Here, the projection of a (generalized) geometric solution of \(\mathcal{M}\) is nothing but an improper affine sphere outside singular points. The singularities of indefinite generalized IA-maps and those of the projection of generalized geometric solutions of \(\mathcal{M}\) share some same properties as mentioned in Remark 3.8. However, there are also different properties between singularities which appear in both formulations. Indeed, we
can find the singularity on indefinite generalized IA-maps which does not appear on the generalized geometric solution of $\mathcal{M}$.

Before proving this, we should review the Ishikawa-Machida’s formulation briefly. In [10], they studied the singularities of graphs $z = f(x, y)$ in $xyz$-space $\mathbb{R}^3$ where $f$ is a solutions of the Monge-Ampère type equation $f_{xx}f_{yy} - f_{xy}^2 = c$. Here, the equation $f_{xx}f_{yy} - f_{xy}^2 = c$ can be considered geometrically in terms of the differential system $\mathcal{M}$ on $xyzpq$-space $\mathbb{R}^5$, which is generated by

\begin{equation}
\omega = cdx \wedge dy - dp \wedge dq,
\end{equation}

where $p, q$ represent $z_x = f_x, z_y = f_y$ respectively. Here, $D = \{ \theta = 0 \}$ is the contact structure on $\mathbb{T}\mathbb{R}^3$. For the graph $z = f(x, y)$ of a solution of the Monge-Ampère type equation $f_{xx}f_{yy} - f_{xy}^2 = c$ in $xyz$-space $\mathbb{R}^3$, we define its lift $L_f : \mathbb{R}^2 \to \mathbb{R}^3$ to $\mathbb{R}^3$ as $L_f(x, y, z) := (x, y, f(x, y), f(x, y), f_y(x, y))$. Obviously, the lift of the graph of $f_{xx}f_{yy} - f_{xy}^2 = c$ annihilates both $\omega$ and $\theta$, that is, it is a Legendrian immersion into $\mathbb{R}^3$ which annihilates $\omega$. Taking this into consideration, a geometric solution (respectively, generalized geometric solution) of $\mathcal{M}$ is defined as a Legendrian immersion with respect to the contact structure $D$ (respectively, a map annihilating $\omega$ which is not necessarily an immersion) of $\mathbb{R}^2$ into $\mathbb{R}^3$, which also annihilates $\theta$.

Among the singularities of improper affine spheres $\psi = (x, y, z) : \mathbb{R}^2 \to \mathbb{R}^3$ which appear as the projection of (generalized) geometric solution $f = (x, y, z, p, q) : \mathbb{R}^2 \to \mathbb{R}^3$ of the Monge-Ampère system $\mathcal{M}$, we are especially interested in singularities of corank 1. In this case, by implicit function theorem, we can take a coordinate system $(u, v)$ of $\mathbb{R}^2$ around $(0, 0)$ such that $\psi(u, v) = (u, y(u, v), z(u, v))$. Since $f^*\theta = 0$, we have

\begin{equation}
z_u = p + qy_u \quad \text{and} \quad z_v = qy_v.
\end{equation}

So the unit normal $\tilde{\nu}$ of $f$ becomes $\tilde{\nu} = \frac{1}{p^2 + q^2 + 1}(-p, -q, 1)$ because $f_u = (1, y_u, z_u)$ and $f_v = (0, y_v, z_v)$. Hence the signed area density of $f$ is $\lambda = (1 + q^2)y_v$, and we can conclude that a point $(u_0, v_0)$ is a singular point if $y_{uv}(u_0, v_0) = 0$ and that a singular point $(u_0, v_0)$ is degenerate (respectively, non-degenerate) if $y_{uvv}(u_0, v_0) = y_{vvv}(u_0, v_0) = 0$ (resp. neither $y_{uvv}(u_0, v_0) \neq 0$ nor $y_{vvv}(u_0, v_0) \neq 0$ holds).

Now, based on the above review, we can prove the following proposition on singularities on improper affine spheres.

**Proposition 6.1.** The germ of the map of Example 5.2 at $(1, 1) \in \mathbb{R}^2$ is not $A$-equivalent to any germ of generalized geometric solution of $\mathcal{M}$.

The proof is based on the following well-known facts about singularities.

**Fact 6.2** (Mather division theorem [3]). Let $F$ be a smooth real-valued function defined on a neighborhood of 0 in $\mathbb{R} \times \mathbb{R}$ such that $F(0, t) = g(t)t^k$ where $g(0) \neq 0$ and $q$ is smooth on some neighborhood of 0 in $\mathbb{R}$. Then given any smooth real-valued function $G$ defined on a neighborhood of 0 in $\mathbb{R} \times \mathbb{R}$, there exist smooth functions $q$ and $r$ such that

\begin{enumerate}
\item $G = qF + r$ on a neighborhood of 0 in $\mathbb{R} \times \mathbb{R}$, and
\item $r(x, t) = \sum_{i = 0}^{k-1} r_i(x)t^i$ for $(t, x) \in \mathbb{R} \times \mathbb{R}$ near 0.
\end{enumerate}

Next, we denote by $E_2$ the set of smooth functions on $\mathbb{R}^2$. For a smooth map germ $f = (f^1, f^2, f^3) : (\mathbb{R}^2, 0) \to (\mathbb{R}^3, 0)$, we denote by $I(f)$ the ideal of $E_2$ generated by $f^1, f^2$ and $f^3$ and define $Q(f) := E_2/I(f)$. 
Fact 6.3 (118). Let $f,g : (\mathbb{R}^2,0) \to (\mathbb{R}^3,0)$ be two map-germs from $\mathbb{R}^2$ to $\mathbb{R}^3$. If $f$ and $g$ are $\mathcal{A}$-equivalent each other, then $Q(f)$ and $Q(g)$ are isomorphic as $\mathbb{R}$-algebras.

Proof of Proposition 6.7. The map of Example 5.2 is concretely expressed as

$$\psi(u,v) = (u^3 + 3uv^2 - u^4 - 6uv^2 - v^4, v^3 + 3u^2v + 4u^3v + 4uv^3, \Phi)$$

where

$$\Phi = \frac{1}{2} u^6 - \frac{3}{2} u^4 v^2 + \frac{3}{2} u^2 v^4 - \frac{1}{2} v^6 - \frac{1}{7} u^7 - 3u^5 v^2 - 5u^3 v^4 - uv^6 - \frac{1}{2} v^8 + 2u^6 v^2 - 3u^4 v^4 + 2u^2 v^6 - \frac{1}{2} v^8.$$ 

Define $\tilde{\psi} := (\tilde{\psi}^1, \tilde{\psi}^2, \tilde{\psi}^3) := \psi(u - 1, v - 1)$, then $Q(\tilde{\psi}) := E_2/(\tilde{\psi}^1, \tilde{\psi}^2, \tilde{\psi}^3)_{E_2} \cong \langle 1, v, v^2 \rangle$. Recall that the map-germ $(\psi, (1, 1))$ (and so $(\tilde{\psi}, (0, 0))$ possesses the following property at a singular point $p \in \Sigma^2$. (Example 5.2)

(i) $p$ is a degenerate singular point.
(ii) $\psi$ is frontal but not front of corank 1 on $p$.
(iii) the set of degenerate singular point around $p$ is locally a smooth curve.

In the following, we will show that for $f = (x, y, z) : \mathbb{R}^2 \to \mathbb{R}^3$, the projection of a (generalized) geometric solution of $\mathcal{M}$ and for $(u_0, v_0) \in \mathbb{R}^2$, if the map-germ $(f, (u_0, v_0))$ is $\mathcal{A}$-equivalent to $(\psi, (0, 0))$ then contradiction occurs. From assumption, $(f, (u_0, v_0))$ also have the above properties (i), (ii) and (iii) because these three conditions are preserved under the same $\mathcal{A}$-equivalent class. By the change of coordinates of $\mathbb{R}^2$ and $\mathbb{R}^3$, we can assume that $(u_0, v_0) = (0, 0)$ and $f(0, 0) = (0, 0, 0)$.

First, the condition (i) and (ii) implies that $x(u, v) = u$ and $y_v(0, 0) = y_{vv}(0, 0) = y_{vvv}(0, 0) = 0$. Therefore the Taylor expansion of $y(u,v)$ around $(0,0)$ becomes as the following form:

$$y(u,v) = a_k v^k q(v) + \sum_{i,j} a_{ij} u^i v^j$$

where $k \geq 3$, $a_k \neq 0$ and $q(v)$ is a smooth function with $q(0) \neq 0$ because otherwise, $y(u,v), z(u,v) \in I(u)$ from (6.2). It follows from (6.2) and (6.3) that $z(u,v) \in I(y(u,v))$. Hence $Q(f) \cong \langle 1, v, \ldots, v^{k-1} \rangle$ and $Q(\tilde{\psi}) \cong Q(f)$ implies that $k = 3$.

By the way, the condition (iii) implies that

(iv) $\{y_v = 0\} \cap \{y_{vv} = 0\}$ is locally a smooth curve.

From (6.3), Fact 5.2 implies that

$$y_v(u,v) = P(u,v) y_{vv}(u,v) + R(u)$$

holds. The condition (iv) implies that $y_v(u,v) = 0$ if $y_{vv}(u,v) = 0$ because $y_{vv}(u,v) = 0$ is locally a smooth curve and that there exists some $v$ for any $u$ such that $y_v(u,v)$ and $y_{vv}(u,v)$ holds around $(0,0)$ unless the singular set is $\{u = 0\}$. Thus $R(u) = 0$ holds for any $u$ around 0. Therefore, locally

$$y_v(u,v) = P(u,v) y_{vv}(u,v)$$

holds. The general solution of (6.4) is

$$y(u,v) = \int \exp \left( \int \frac{dv}{P(u,v)} + \alpha(u) \right) dv + \beta(u),$$

but this contradicts to the condition (iv) because $\{y_v = 0\} = \emptyset$. 

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