Some nontrivial secondary adams differentials on the fourth line

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Abstract. Let $p \geq 5$ be an odd prime. Using the correspondence between secondary Admas differentials and secondary algebraic Novikov differentials, we compute four families of nontrivial secondary differentials on the fourth line of the Adams spectral sequence. We also recover all secondary differentials on the first three lines of the Adams spectral sequence.

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1. Introduction

The Adams spectral sequence (ASS) is one of the most useful tools to compute the stable homotopy groups of the sphere $\pi_*(S)$. The ASS has $E_2$-page $Ext_{A_*}^{p,n}(F_p, F_p)$, where $A_*$ is the dual mod $p$ Steenrod algebra.

In this paper, we always assume $p$ is an odd prime. Then we have

$$A_* = P(\xi_1, \xi_2, \cdots) \otimes E[\tau_0, \tau_1, \tau_2, \cdots],$$

where $P(\xi_1, \xi_2, \cdots)$ is a polynomial algebra with coefficients in $F_p$, and $E[\tau_0, \tau_1, \tau_2, \cdots]$ is an exterior algebra with coefficients in $F_p$.

The Adams-Novikov spectral sequence (ANSS) is another useful tool for computing $\pi_*(S)$. The ANSS has $E_2$-page $Ext_{BP_*BP}^{p,n}(BP_*, BP_*)$, where $BP$ denotes
the Brown-Peterson spectrum. We have

\[ BP_* := \pi_*(BP) = \mathbb{Z}(p)[v_1, v_2, \cdots] , \quad BP_*BP = BP_*[t_1, t_2, \cdots] \]

where \( \mathbb{Z}(p) \) denotes the integers localized at \( p \).

The Adams-Novikov \( E_2 \)-page can be computed via the algebraic Novikov spectral sequence (algNSS) \([10, 14]\). The \( E_2 \)-page of the algNSS has the form \( Ext^{s, t}_{\hat{p}}(\mathbb{F}_p, I^k/I^{k+1}) \), where \( I \) denotes the ideal \( (p, v_1, v_2, \cdots) \subset BP_* \), and \( P_* = BP_*BP/I = P[t_1, t_2, \cdots] \) is the \( \mathbb{F}_p \)-coefficient polynomial algebra. Here, we have re-indexed the pages to align with the notations in Gheorghe-Wang-Xu \([4]\) and Isaksen-Wang-Xu \([6]\).

The \( E_2 \)-page of the Adams spectral sequence can also be computed via another spectral sequence, called the Cartan-Eilenberg spectral sequence (CESS) \([3, 15]\). For odd prime \( p \), the \( E_2 \)-page of the CESS coincides with the \( E_2 \)-page of the algNSS. Then, we have the following diagram of spectral sequences.

\[
\begin{array}{ccc}
Ext^{s, t}_{\hat{p}}(\mathbb{F}_p, I^k/I^{k+1}) & \xrightarrow{CESS} & Ext^{s+k+k, t+k}_{A_*}(\mathbb{F}_p, \mathbb{F}_p) \\
\downarrow \text{algNSS} & & \downarrow \text{ASS} \\
Ext^{s, t}_{BP_*BP}(BP_*, BP_*) & \xrightarrow{ANSS} & \pi_{t-s}(S)
\end{array}
\]

In practice, the main difficulty of computing with the ASS is that the Adams differentials \( d_2^{\text{Adams}} \)'s are difficult to be determined in general. On the other hand, the algebraic Novikov differentials \( d_2^{\text{alg}} \)'s are much easier to be computed. This is because the entire construction of the algNSS is purely algebraic. Computing \( d_2^{\text{alg}} \)'s does not require any topological background knowledge. It turns out that when \( r = 2 \), there is a direct correspondence between \( d_2^{\text{Adams}} \)'s and \( d_2^{\text{alg}} \)'s.

**Theorem 1.1** (Novikov \([14]\), Andrews-Miller \([2, 11]\)). Let \( z \in Ext^{s+k+k, t+k}_{A_*}(\mathbb{F}_p, \mathbb{F}_p) \) be a nontrivial element detected in the CESS by \( x \in Ext^{s, t}_{\hat{p}}(\mathbb{F}_p, I^k/I^{k+1}) \). Regard \( x \) as an element in the algNSS, then the secondary algebraic Novikov differential \( d_2^{\text{alg}}(x) \) represents the secondary Adams differential \( d_2^{\text{Adams}}(z) \).

Let \( p \geq 5 \). A complete list of generators together with their \( d_2^{\text{Adams}}(z) \) has been determined for the first three lines of the Adams \( E_2 \)-page, i.e. \( Ext^{s, t}_{A_*}(\mathbb{F}_p, \mathbb{F}_p) \) with \( s = 1, 2, 3 \) (see \([1, 7, 12, 16, 17, 18]\)). Meanwhile, only partial results are known for the fourth line \( Ext^{4, s}_{A_*}(\mathbb{F}_p, \mathbb{F}_p) \) (see, for example, \([19]\)).

In this paper, we demonstrate a practical computing strategy to determine \( d_2^{\text{Adams}} \)'s by computing their corresponding \( d_2^{\text{alg}} \)'s. We will work on several explicit examples and provide detailed proof. Our main result is the following.
Theorem 4.4. There are nontrivial secondary Adams differentials given as follows:

1. \( \partial^2_2(h_{4,i}h_{3,j}g_i) = a_0b_{4,i-1}h_{3,j}g_i, \text{ for } i \geq 1. \)
2. \( \partial^2_2(h_{4,i}h_{3,j+1}k_{i+2}) = a_0b_{4,i-1}h_{3,j+1}k_{i+2}, \text{ for } i \geq 1. \)
3. \( \partial^2_2(h_{3,j}h_{2,i}h_{i+1}k) = a_0b_{3,i-1}h_{2,i}h_{i+1}k, \text{ for } i \geq 1. \)

Remark 1.2. Here, we follow the conventions of [17, 19] to name Adams-2-page elements by their May spectral sequence (MSS) representatives, compare with Table 2 and Table 3. We would like to comment more explicitly on the indeterminacy of these classes. For example, the result of (1) should be interpreted as follows. If an element \( h_{4,i}h_{3,j}g_i \) has MSS representative \( h_{4,i}h_{3,j}h_{1,i} \), then its secondary Adams differential \( \partial^2_2(h_{4,i}h_{3,j}g_i) \) has MSS representative \( a_0b_{4,i-1}h_{3,j}g_i \). More details of the MSS are reviewed in Section 3.

It is straightforward to verify that these four families of elements are indecomposable, i.e., they can not be written as products of elements from the first three lines. Consequently, one can not deduce the differentials simply via Leibniz rule.

From our point of view, the practical computational strategy here is possibly more interesting than the result itself. To further demonstrate this, in Section 5, we use the same strategy to recover all secondary Adams differentials on the first three lines.

Previously, the nontrivial Adams differentials on the third line were computed in [17] using matrix Massey products [9]. Comparatively, our computation has the following advantages: (i) Our computations can be easily adapted to analyze other \( \partial^2_2 \)'s of interest. On the contrary, the matrix Massey product method could fail when the relevant indeterminacy is nontrivial; (ii) Our computations of the algebraic Novikov differentials are routine and purely algebraic. Such computations are comparatively more straightforward than the previous ones using matrix Massey products.

Organization of the paper. In Section 2, we review the algebraic structures and constructions related to Hopf algebroids. These structures are fundamental to later computations. In Section 3, we discuss several spectral sequences we use in this paper, including the algNSS, the CESS, and the May spectral sequence. In Section 4, we compute relevant algebraic Novikov differentials and prove Theorem 4.4. In Section 5, we use the same computational strategy to recover the secondary Adams differentials on the first three lines.

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2. Hopf algebroids

In this section, we review the definition as well as two important examples of Hopf algebroids. We will also recall the associated cobar complex construction.

**Definition 2.1** ([15] Definition A1.1.1). A Hopf algebroid over a commutative ring $K$ is a pair $(A, \Gamma)$ of commutative $K$-algebras with the following structure maps

- left unit map $\eta_L : A \to \Gamma$
- right unit map $\eta_R : A \to \Gamma$
- coproduct map $\Delta : \Gamma \to \Gamma \otimes_A \Gamma$
- counit map $\varepsilon : \Gamma \to A$
- conjugation map $c : \Gamma \to \Gamma$

such that for any other commutative $K$-algebra $B$, the two sets of $K$-morphisms, $\text{Hom}_K(A, B)$ and $\text{Hom}_K(\Gamma, B)$, are the objects and morphisms of a groupoid.

**2.1. The Hopf algebroid $(BP_*, BP_*BP)$.** An important example of Hopf algebroids is $(BP_*, BP_*BP)$ [5, 12, 15]. Recall that we have

$$BP_* := \pi_*(BP) = \mathbb{Z}[[v_1, v_2, \cdots]], \quad BP_*BP = BP_*[t_1, t_2, \cdots]$$

(1)

We also have

$$H_*(BP) = \mathbb{Z}[[m_1, m_2, \cdots]]$$

(2)

where $|v_n| = |t_n| = |m_n| = 2(p^n - 1)$.

**Notations 2.2.** Throughout this paper, we denote $v_0 = p$, and $m_0 = t_0 = 1$.

The Hurewicz map induces an embedding

$$i : BP_* \to H_*(BP)$$

$$v_n \mapsto pm_n - \sum_{i=1}^{n-1} v_i^{p^i} m_i$$

(3)

We can describe the structure maps of the Hopf algebroid $(BP_*, BP_*BP)$ as follows.

The left unit and right unit maps $\eta_L, \eta_R : BP_* \to BP_*BP$ are determined by

$$\eta_L(v_n) = v_n$$

(4)

$$\eta_R(m_n) = \sum_{i+j=n} m_it_j^{p^i}$$

(5)

The coproduct map $\Delta : BP_*BP \to BP_*BP \otimes_{BP_*} BP_*BP$ is determined by

$$\sum_{i+j=n} m_i(\Delta t_j)^{p^i} = \sum_{i+j+k=n} m_it_j^{p^i} \otimes t_k^{p^{i+j}}$$

(6)

The counit map $\varepsilon : BP_*BP \to BP_*$ is determined by

$$\varepsilon(v_n) = v_n, \quad \varepsilon(t_n) = 0.$$  

(7)
The conjugation map $c : BP_*BP \to BP_*BP$ is determined by

$$
\sum_{i+j+k=n} m_i t_j^{p^i} c(t_k^{p^j}) = m_n. \quad (8)
$$

In practice, it is more convenient to work with $\eta_R(v_n)$ instead of $\eta_R(m_n)$.

Let $I$ denote the ideal $(p, v_1, v_2, \cdots) \subset BP_*$. Then $I$ is an invariant ideal as a $BP_*BP$-comodule, in other words, we have $\eta_L(I) \cdot BP_*BP = BP_*BP \cdot \eta_R(I)$. For $k \geq 0$, we let $I^k \cdot BP_*BP$ denote $\eta_L(I^k) \cdot BP_*BP = BP_*BP \cdot \eta_R(I^k)$.

We have the following formulas.

**Proposition 2.3.** Let $n \geq 0$. The right unit map $\eta_R : BP_* \to BP_*BP$ satisfies

$$
\eta_R(v_n) \equiv \sum_{i=0}^{n} v_i t_{n-i}^{p^i} \mod I^p \cdot BP_*BP. \quad (9)
$$

**Proof.** We prove by induction on $n$. The case for $n = 0$ is trivial. Now suppose (9) is true for $0 \leq i \leq n - 1$. Then, in particular, $\eta_R(v_i) \in I \cdot BP_*BP$ for $0 \leq i \leq n - 1$. Note (3) implies

$$
v_n \equiv pm_n \mod I^p H_s(BP). \quad (10)
$$

for $n \geq 0$. Then, direct computation shows

$$
\eta_R(v_n) = p \eta_R(m_n) - \sum_{i=1}^{n-1} \eta_R(v_{n-i}) p^i \eta_R(m_i) \quad \text{ (by (3))}
$$

$$
\equiv p \eta_R(m_n) \mod I^p \cdot BP_*BP \quad (11)
$$

$$
\equiv p \sum_{i=0}^{n} m_i t_{n-i}^{p^i} \mod I^p \cdot BP_*BP \quad \text{ (by (5))}
$$

$$
\equiv \sum_{i=0}^{n} v_i t_{n-i}^{p^i} \mod I^p \cdot BP_*BP \quad \text{ (by (10))}
$$

\[ \square \]

Similarly, we could obtain the following formulas for $\Delta(t_n)$.

**Proposition 2.4.** For $n \geq 0$, we have

$$
\Delta(t_n) = \sum_{k=0}^{n} t_{n-k} \otimes t_k^{p^{n-k}} - \sum_{i=1}^{n-1} v_i b_{n-i-1} \mod I^2 \cdot BP_*BP \otimes_{BP_*BP} BP_*BP \quad (12)
$$

where we denote

$$
b_{i,j} = \frac{1}{p} \left[ \left( \sum_{k=0}^{i} t_{i-k} \otimes t_k^{p^{i-k}} \right) t_{i+1}^{p+1} - \sum_{k=0}^{i} t_{i-k} \otimes t_k^{p^{i-k+1}} \right] \quad (13)
$$

for $i \geq 1, j \geq 0$. 
Proof. We prove by induction on $n$. The case for $n = 0$ is trivial. Now suppose (12) is true for $0 \leq i \leq n - 1$. Then, direct computation shows

$$
\Delta(t_n) = \sum_{i+j=n} m_i(\Delta t_j)^{p^i} - \sum_{i=1}^n m_i(\Delta t_{n-i})^{p^i}
$$

$$
= \sum_{i+j+k=n} m_i t_j^{p^i} \otimes t_k^{p^{i+j}} - \sum_{i=1}^n m_i(\Delta t_{n-i})^{p^i}
$$

$$
= \sum_{k=0}^{n} t_{n-k} \otimes t_k^{p^{n-k}} + \sum_{i=1}^n m_i(\sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-i-k}}) - \sum_{i=1}^n m_i(\Delta t_{n-i})^{p^i}
$$

$$
= \sum_{k=0}^{n} t_{n-k} \otimes t_k^{p^{n-k}} - \sum_{i=1}^n m_i[(\Delta t_{n-i})^{p^i} - \sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-k}}]
$$

Modulo $I^2 \cdot BP_*BP \otimes_{BP_*BP} BP_*BP$, we have

$$
\sum_{i=1}^n m_i[(\Delta t_{n-i})^{p^i} - \sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-i-k}}]
$$

$$
= \sum_{i=1}^n m_i[(\sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-i-k}})^{p^i} - \sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-i-k}}]
$$

$$
= \sum_{i=1}^{n-1} p m_i \cdot \frac{1}{p}[(\sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-i-k}})^{p^i} - \sum_{k=0}^{n-i} t_{n-i-k} \otimes t_k^{p^{n-i-k}}]
$$

$$
= \sum_{i=1}^{n-1} v_i b_{n-i-1}
$$

This completes the proof. \hfill \square

2.2. The dual Steenrod algebra $A_*$. The Steenrod algebra provides another important example of Hopf algebroids.

Let $A_*$ denote the dual mod $p$ Steenrod algebra for an odd prime $p$, we have \cite{13}

$$
A_* = P[\xi_1, \xi_2, \cdots] \otimes E[\tau_0, \tau_1, \tau_2, \cdots]
$$

as an algebra, where $P[\xi_1, \xi_2, \cdots]$ is a polynomial algebra with coefficients in $\mathbb{F}_p$, $E[\tau_0, \tau_1, \tau_2, \cdots]$ is an exterior algebra with coefficients in $\mathbb{F}_p$. For the internal degrees, we have $|\xi_n| = 2(p^n - 1)$, $|\tau_n| = 2p^n - 1$. We also denote $\xi_0 = 1$.

One can show $A_*$ is a Hopf algebra over $\mathbb{F}_p$. In particular, $(\mathbb{F}_p, A_*)$ has a Hopf algebroid structure. We can describe the structure maps as follows [13].

The left unit $\eta_L : \mathbb{F}_p \to A_*$, right unit $\eta_R : \mathbb{F}_p \to A_*$, and counit $\epsilon : A_* \to \mathbb{F}_p$ maps are all isomorphisms in dimension 0.
On generators, the coproduct $\Delta : A_* \to A_* \otimes A_*$ is given by:

$$\Delta \xi_n = \sum_{i=0}^{n} \xi_{n-i}^{p^i} \otimes \xi_i, \quad \Delta \tau_n = \tau_n \otimes 1 + \sum_{i=0}^{n} \xi_{n-i}^{p^i} \otimes \tau_i$$ (16)

The conjugation map $c : A_* \to A_*$ is an algebra map given recursively by

$$c(\xi_0) = 1, \quad \sum_{i=0}^{n} \xi_{n-i}^{p^i} c(\xi_i) = 0, n > 0,$$

$$\tau_n + \sum_{i=0}^{n} \xi_{n-i}^{p^i} c(\tau_i) = 0, n \geq 0.$$ (17) (18)

For our computational purposes, we prefer to use a different set of generators. We denote $t_n = c(\xi_n), n \geq 1$, and $\tau_n = c(\tau_n), n \geq 0$. We also denote $t_0 = 1$.

**Proposition 2.5.** Let $p$ be an odd prime, we can write

$$A_* = P[t_1, t_2, \cdots] \otimes E[\tilde{\tau}_0, \tilde{\tau}_1, \tilde{\tau}_2, \cdots]$$ (19)

as an algebra, where $|t_n| = 2(p^n - 1), |\tilde{\tau}_n| = 2p^n - 1$. Moreover, the coproduct $\Delta : A_* \to A_* \otimes A_*$ is given by:

$$\Delta t_n = \sum_{i=0}^{n} t_i \otimes t_{n-i}^{p^i}, \quad \Delta \tilde{\tau}_n = \sum_{i=0}^{n} \tilde{\tau}_i \otimes t_{n-i}^{p^i} + 1 \otimes \tilde{\tau}_n$$ (20)

**Proof.** It is straightforward to deduce the coproduct formulas by induction on $n$. Here, we outline the strategy to prove (20) for $t_n$. The formula for $\tilde{\tau}_n$ can be verified similarly.

The case for $n = 0$ is trivial. Now, suppose (20) is true for $0 \leq m \leq n - 1$. Note (17) implies

$$\sum_{i=0}^{n-1} (\Delta \xi_{n-i})^{p^i} (\Delta t_i) + \Delta t_n = 0$$

To deduce the desired result, it suffices to show

$$\sum_{i=0}^{n-1} (\Delta \xi_{n-i})^{p^i} (\Delta t_i) + \sum_{i=0}^{n} t_i \otimes t_{n-i}^{p^i} = 0$$
Indeed, we have
\[
\sum_{i=0}^{n-1} (\Delta \xi_{n-i})^p (\Delta t_i) + \sum_{i=0}^{n} t_i \otimes t_{n-i}^p
\]
\[
= \sum_{i=0}^{n} \sum_{j=0}^{n-i} \xi_{n-i-j}^p \otimes \xi_j^p (\sum_{k=0}^{i} t_k \otimes t_{i-k}^p)
\]
\[
= \sum_{i=0}^{n} \sum_{j=0}^{n-i} \xi_{n-i-j}^p \otimes \xi_j^p (\sum_{k=0}^{i} t_k \otimes t_{i-k}^p)
\]
\[
= \sum_{j+k+\ell=n} \xi_j^p \otimes \xi_{\ell}^p \otimes \xi_0^p t_k \otimes t_{\ell}^p t_{\ell}^p
\]
\[
= \sum_{r+k<n} \xi_r^p \otimes (\sum_{j=s=n-k-r} \xi_j^p t_s^p + \sum_{r+k=n} \xi_r^p) t_k \otimes 1
\]
\[
= 0
\]
\[\square\]

**Remark 2.6.** The advantage of using the new set of generators is that, as we will see in Section 3, $c(\xi_n)$ corresponds to the generator $t_n \in BP_sBP$ and $c(t_n)$ corresponds to $u_n \in BP_v$. Hence, we abuse the notation and denote $c(\xi_n)$ as $t_n$ when no confusion arises.

### 2.3. Cobar complexes.

**Definition 2.7.** Let $(A, \Gamma)$ be a Hopf algebroid. A **right $\Gamma$-comodule** $M$ is a right $A$-module $M$ together with a right $A$-linear map $\psi : M \to M \otimes_A \Gamma$ which is counitary and coassociative, i.e., the following diagrams commute.

\[
\begin{array}{ccc}
M & \xrightarrow{\psi} & M \otimes_A \Gamma \\
\downarrow & & \downarrow \\
M & \xrightarrow{\coprod} & M \otimes_A \Delta \\
\end{array}
\]

Left $\Gamma$-comodules are defined similarly.

**Definition 2.8.** Let $(A, \Gamma)$ be a Hopf algebroid. Let $M$ be a right $\Gamma$-comodule. The **cobar complex** $\Omega^{s \times r}_\Gamma(M)$ is a cochain complex with
\[
\Omega^{s \times r}_\Gamma(M) = M \otimes_A \Gamma^{\otimes s}
\]
where $\Gamma$ is the augmentation ideal of $\epsilon : \Gamma \to A$. The differentials $d : \Omega^{s \times r}_\Gamma(M) \to \Omega^{s+1 \times r}_\Gamma(M)$ are given by
\[
d(m \otimes x_1 \otimes x_2 \otimes \cdots \otimes x_n) = - (\psi(m) - m \otimes 1) \otimes x_1 \otimes x_2 \otimes \cdots \otimes x_n
\]
\[
- \sum_{i=1}^{s} (-1)^{y_{j_i}} m \otimes x_1 \otimes \cdots \otimes x_{i-1} \otimes (\sum_{j_i} x'_{j_i} \otimes x''_{j_i}) \otimes x_{i+1} \otimes \cdots \otimes x_n
\]
where
\[ \sum_{j_i} x_{i,j_i}' \otimes x_{i,j_i}'' = \Delta(x_i) - 1 \otimes x_i - x_i \otimes 1 \]
\[ \lambda_{i,j_i} = i + |x_1| + \cdots + |x_{i-1}| + |x_i'| \]

The cohomology of \( \Omega^s_{A} (M) \) is \( \text{Ext}^s_A (M, M) \) (see [15, Section A1.2]).

3. Some relevant spectral sequences

In this section, we review the construction and properties of some relevant spectral sequences, including the algebraic Novikov spectral sequence (algNSS), the Cartan-Eilenberg spectral sequence (CESS), and the May spectral sequence (MSS). These spectral sequences will be used in later computations.

3.1. The algebraic Novikov spectral sequence. Let \( I \) be the ideal of \( BP_* \) generated by \( (p, v_1, v_2, \cdots) \). The ideal \( I \) induces a filtration
\[ BP_* = I^0 \supset I^1 \supset I^2 \supset I^3 \supset \cdots \supset I^k \supset I^{k+1} \supset \cdots \] (22)

Consider \( y = a p^{k_0} v_1^{k_1} v_2^{k_2} \cdots \in BP_* \), where \( a \in \mathbb{Z} \) is invertible. We let \( l(y) = \Sigma k_i \) denote the length of \( y \). Then \( y \in I^k \) if and only if \( l(y) \geq k \).

Let \( E^0_*BP_* \) denote the associated graded object, where \( E^0_*BP_* := I^k / I^{k+1} \).
We have
\[ E^0_*BP_* = \bigoplus_{k \geq 0} I^k / I^{k+1} = \mathbb{F}_p [q_0, q_1, q_2, \cdots] \] (23)
is a \( \mathbb{F}_p \)-coefficient polynomial algebra, where the generator \( q_i \) corresponds to \( v_i \). \( I^k / I^{k+1} \) corresponds to those homogeneous polynomials of degree \( k \).

Similarly, we can filter \( BP_*BP \). Denote
\[ F^kBP_*BP := \eta_l(I^k)BP_*BP = BP_*BP \eta_k(I^k) \]
We define the associated graded object \( E^k_*BP_*BP := F^kBP_*BP / F^{k+1}BP_*BP \).
The filtration of \( BP_* \) and \( BP_*BP \) together induces a filtration on \( \Omega_{BP_*BP}(BP_*). \) Such filtration induces an associated spectral sequence [15, A1.3.9] converging to \( \text{Ext}^s_{BP_*BP}(BP_*BP) \).

**Theorem 3.1** ([10, 14]). There is a spectral sequence, called the algebraic Novikov spectral sequence (algNSS), converging to \( \text{Ext}^s_{BP_*BP}(BP_*BP) \) with \( E_2 \)-page
\[ E_2^{s,t,k} = \text{Ext}^s_{BP_*BP}(\mathbb{F}_p, I^k / I^{k+1}) \]
and \( d_r^{\text{alg}} : E_r^{s,t,k} \to E_r^{s+1,t,k+r-1} \), where
\[ P_* := E_0BP_*BP \otimes_{E_0BP_*BP} \mathbb{F}_p = BP_*BP / I = P[t_1, t_2, \cdots] \] (24)
is the \( \mathbb{F}_p \)-coefficient polynomial algebra.

**Remark 3.2.** Our index of pages here is different from the ones used in [2, 15]. We have re-indexed the spectral sequence to align with the notations in [4, 6].
3.2. The Cartan-Eilenberg spectral sequence. Let \( A_* \) denote the dual Steenrod algebra for an odd prime \( p \). Recall from Proposition 2.5 that we have
\[
A_* = P[t_1, t_2, \cdots] \otimes E[\bar{\tau}_0, \bar{\tau}_1, \bar{\tau}_2, \cdots]
\]

Let \( P_* \) denote \( P[t_1, t_2, \cdots] \subset A_* \). Let \( E_* \) denote \( E[\bar{\tau}_0, \bar{\tau}_1, \bar{\tau}_2, \cdots] \). Then
\[
P_* \to A_* \to E_*
\]
is an extension of Hopf algebras \([15, A1.1.15]\), which induces a spectral sequence \([15, A1.3.14]\) converging to \( \text{Ext}^{*,*}(F_p, F_p) \).

**Theorem 3.3** ([15] Theorem 4.4.3, 4.4.4). Let \( p \) be an odd prime. There is a spectral sequence, called the Cartan-Eilenberg spectral sequence (CESS), converging to \( \text{Ext}^{*,*}(F_p, F_p) \) with \( E_2 \)-page
\[
E_2^{s,t} = \text{Ext}_{E_*}^{s,t}(F_p, \text{Ext}_{E_*}^t(F_p, F_p))
\]
and \( d_r : E_r^{s,t} \to E_r^{s+r, t-s-r+1} \). Moreover, one can prove the following results:
(a) \( \text{Ext}_{E_*}^{*,*}(F_p, F_p) = P[a_0, a_1, \cdots] \) is a polynomial algebra with generator \( a_i \in \text{Ext}^{1,2p^i-1} \) represented in the associated cobar complex \( \Omega_{E_*}(F_p) \) by \( [\bar{\tau}_i] \).
(b) The \( P_* \)-coaction on \( \text{Ext}_{E_*}(F_p, F_p) \) is given by
\[
\psi(a_n) = \sum_{i=0}^{n} a_i \otimes t_n^{p^i}
\]
(c) The CESS collapses from \( E_2 \) with no nontrivial extensions.
(d) There is an isomorphism
\[
\text{Ext}_{E_*}^t(F_p, I^k/I^{k+1}) \cong \text{Ext}_{E_*}^{t+k}(F_p, \text{Ext}_{E_*}^k(F_p, F_p))
\]
between the \( E_2 \)-page of the algNSS and the \( E_2 \)-page of the CESS.

The (d) part shows the two \( \text{Ext} \) groups are isomorphic (up to degree shifting). Moreover, we can show the two associated cobar complexes are isomorphic (up to a shifting of degrees). More precisely, there is a natural isomorphism
\[
\Omega_{P_*}(I^k/I^{k+1}) \cong \Omega_{P_*}(\text{Ext}_{E_*}^k(F_p, F_p))
\]
sending \( t_i \) to \( t_i \) and \( q_l \) to \( a_i \).

Indeed, by Theorem 3.3 (a), \( \text{Ext}_{E_*}^k(F_p, F_p) \) is the homogeneous degree \( k \) part of the polynomial \( P[a_0, a_1, a_2, \cdots] \). Hence \( I^k/I^{k+1} \cong \text{Ext}_{E_*}^k(F_p, F_p) \). If \( x \in I^k/I^{k+1} \) has inner degree \( t \), then its corresponding element \( \bar{x} \in \text{Ext}_{E_*}^k(F_p, F_p) \) has inner degree \( t + k \). This degree shifting is a consequence of the fact that \( |q_l| = 2(p^l - 1) = |a_i| - 1 \). Moreover, the comodule structure map \( \psi : I^k/I^{k+1} \to I^k/I^{k+1} \otimes P_* \) induced from (9) is given by
\[
\psi(q_n) = \sum_{i=0}^{n} q_i \otimes t_n^{p^i}
\]
which also agrees with (25).
Notations 3.4. In this paper, we often refer to $E_2$-terms of the algNSS by their representative in the cobar complex $\Omega_\mathcal{F}^s(I^k/I^{k+1})$. For example, we let $q_0 \otimes t^p_1$ denote its homology class in $\text{Ext}^1_{\mathcal{F}}(\mathbb{F}_p, I/I^2)$. The correspondence between different $E_2$-pages becomes clear under this naming convention. For example, $q_0 \otimes t^p_1 \in \text{Ext}^1_{\mathcal{F}}(\mathbb{F}_p, I/I^2)$ in the algNSS corresponds to

$$a_0 \otimes t^p_1 \in \text{Ext}^1_{\mathcal{F}}(\mathbb{F}_p, \text{Ext}^1_{E_s}(\mathbb{F}_p, \mathbb{F}_p))$$

in the CESS, which represents $\tilde{r}_0 \otimes t^p_1 \in \text{Ext}^2_{\mathcal{A}_s}(\mathbb{F}_p, \mathbb{F}_p)$ in the ASS.

3.3. The May spectral sequence. The $E_2$-terms $\text{Ext}^s_{\mathcal{A}_s}(\mathbb{F}_p, \mathbb{F}_p)$ of the Adams spectral sequence could be computed via the cobar complex $\Omega^s_{\mathcal{A}_s}(\mathbb{F}_p)$. In practice, we could simplify such computations by filtering $\Omega^s_{\mathcal{A}_s}(\mathbb{F}_p)$.

Theorem 3.5 ([8], [15] Theorem 3.2.5). Let $p$ be an odd prime. $\mathcal{A}_s$ can be given an increasing filtration by setting the May degree $M(t^p_i) = M(\tilde{r}_{i-1}) = 2i - 1$ for $i \geq 0$. The filtration of $\mathcal{A}_s$ naturally induces a filtration of $\Omega^s_{\mathcal{A}_s}(\mathbb{F}_p)$. The associated spectral sequence converging to $\text{Ext}^s_{\mathcal{A}_s}(\mathbb{F}_p, \mathbb{F}_p)$ is called the May spectral sequence (MSS). The MSS has $E_1$ page

$$E_1^{s,t,M} = E[h_{i,j} | i \geq 1, j \geq 0] \otimes P[b_{i,j} | i \geq 1, j \geq 0] \otimes P[a_i | i \geq 0]$$

(29)

and $d_r : E_r^{s,t,M} \to E_r^{s+1,t,M-r}$, where

$$h_{i,j} = [t^p_i] \in E_1^{1,2(p^i-1)p^j, 2i-1}$$

$$b_{i,j} = \left[ \sum_{k=1}^{p-1} \binom{p}{k} \frac{p}{k} (t^p_i)^k \otimes (t^p_i)^{p-k} \right] \in E_1^{2,2(p^i-1)p^j, p(2i-1)}$$

$$a_i = [\tilde{r}_i] \in E_1^{1,2p^i-1,2i+1}$$

(30)

Remark 3.6. Technically, we could denote the generator by $b_{i,j}$ instead of $b_{i,j}$ to avoid possible confusion with the element

$$b_{i,j} = \frac{1}{p} \left[ \left( \sum_{k=0}^{i} t_{i-k} \otimes t_{k}^{p^{i-k}} \right)^{p+j+1} - \sum_{k=0}^{i} t_{i-k}^{p+j+1} \otimes t_{k}^{p^{i-k+j+1}} \right]$$

defined in (13). However, let $x$ be the element in $\Omega^s_{\mathcal{A}_s}(\mathbb{F}_p)$ corresponding to $b_{i,j}$ (Notations 3.4). Note $\Omega^s_{\mathcal{A}_s}(\mathbb{F}_p)$ has coefficient $\mathbb{F}_p$, we have

$$x = \frac{1}{p} \sum_{k_1 \neq k_2} \sum_{i=1}^{p-1} \binom{p}{t^p_i} (t_{i-k_1} \otimes t_{k_1}^{p^{i-k_1}}) t^p_i (t_{i-k_2} \otimes t_{k_2}^{p^{i-k_2}}) (p-t)^p$$

$$= \frac{1}{p} \sum_{k_1 \neq k_2} \sum_{i=1}^{p-1} \binom{p}{t^p_i} (t_{i-k_1} t_{i-k_2}^{p^{i-k_1}} (p-t)^p) t_{k_1}^{p^{i-k_1}} t_{k_2}^{p^{i-k_2}}$$

$$+ \frac{1}{p} \left( \sum_{k=0}^{i} t_{i-k} \otimes t_{k}^{p^{i-k}} \right)^{p+j+1} - \sum_{k=0}^{i} t_{i-k}^{p+j+1} \otimes t_{k}^{p^{i-k+j+1}}$$
Its May filtration leading term is

\[ \frac{1}{p} \sum_{i=1}^{b-1} \binom{p}{i} t_i t_i^{(p-t)p} = \tilde{b}_{i,j} \]

Therefore, we often abuse the notation and also denote $\tilde{b}_{i,j}$ by $b_{i,j}$.

Note that we can analogously define an increasing filtration on $\Omega_{E_2}(t^k / I^{k+1})$ (hence also on $\Omega_{E_2}(\text{Ext}_E^k (\mathbb{F}_p, \mathbb{F}_p))$ by setting the May degree

\[ M(t_i^{p^j}) = M(q_{i-1}) = 2i - 1 \]

for $i - 1, j \geq 0$. We observe the following structure maps:

\[ \psi(q_n) = \sum_{i=0}^{n} q_i \otimes t_i^{p^j}, \quad \Delta t_n = \sum_{i=0}^{n} t_i \otimes t_i^{p^j}. \quad (31) \]

For $i < n$, we have $M(q_n) = 2n + 1 \geq 2n = 2i + 1 + 2(n - i) - 1 = M(q_i \otimes t_i^{p^j})$.

Similarly, for $0 < i < n$, $M(t_n) = 2n - 1 \geq 2n - 2 = M(t_i \otimes t_i^{p^j})$. Let $d$ denote the differential of the cobar complex $\Omega_{E_2}(t^k / I^{k+1})$ (see Definition 2.8). Then $d$ respects this May filtration. Hence, we can talk about the May filtration of the algNSS $E_2$-terms. Moreover, the May filtration of the elements in the algNSS $E_2$-page agrees with the May filtration of the corresponding elements in the ASS $E_2$-page (see Notations 3.4).

### 4. Secondary Adams differentials on the fourth line

In this section, we prove our main result Theorem 4.4. Using Theorem 1.1, we determine these secondary Adams differentials $d^{Adams}_2$ by computing their corresponding secondary algebraic Novikov differentials $d^{alg}_2$.

Our computational strategy in this paper can be summarized as follows:

1. Let $x$ be an element in the Adams $E_2$-page. Let $l$ be the MSS representative of $x$.
2. As stated in Notations 3.4, we find the the element $x'$ (resp. $l'$) in the algebraic Novikov spectral sequence corresponding to $x$ (resp. $l$). We deduce $l'$ is the May filtration leading term of $x'$.
3. Through a careful analysis of $l'$, we determine the May filtration leading term $y'$ of $d^{alg}_2(x')$.
4. Let $y$ be the element in the MSS corresponding to $y'$. Then we conclude $d^{Adams}_2(x) = y$.

In particular, we will use Table 1 for the four families of Adams $E_2$-terms in Theorem 4.4.

Now we start the actual computations.
Lemma 4.1. Let \( d \) denote the differential in the cobar complex \( \Omega^*_{BP, BP}(BP_*) \) (Definition 2.8). Let \( n, i \geq 1 \), we have

\[
d(t^p_n) = \sum_{k=1}^{n-1} t^p_{n-k} \otimes t^p_k + p b_{n,i-1} \mod I^2 \cdot BP_* BP \otimes_{BP_* BP} BP_* BP
\]

Proof. After reduction module \( I^2 \cdot BP_* BP \otimes_{BP_* BP} BP_* BP \), we have

\[
d(t^p_n) = \Delta(t^p_n) - 1 \otimes t^p_n - t^p_n \otimes 1
\]

\[
= \sum_{k=0}^{n} t_{n-k} \otimes t^p_k - \sum_{i=1}^{n-1} v_i b_{n-i,i-1} t^p_i - 1 \otimes t^p_n - t^p_n \otimes 1 \text{ (by (12))}
\]

\[
= \sum_{k=0}^{n} t_{n-k} \otimes t^p_k - 1 \otimes t^p_n - t^p_n \otimes 1
\]

\[
= \sum_{k=1}^{n-1} t^p_{n-k} \otimes t^p_k + p b_{n,i-1} \text{ (compare with (13))}
\]

\[\square\]

Proposition 4.2. Let \( x \in \text{Ext}_{P_*}^4(F_p, BP_*/I) \) be an element in the E_2-page of the algNSS such that \( x \) has May filtration leading term \( t_4^p \otimes t_2^p \otimes t_2^p \otimes t_1^p \), where \( i \geq 1 \). Then \( d^\text{alg}_2(x) \) has May filtration leading term \( q_0 b_{4,i-1} \otimes t_3^p \otimes t_2^p \otimes t_1^p \).

Proof. We will compute \( d^\text{alg}_2(x) \) as follows. First, we will find a representative \( \hat{x} \) of \( x \) in \( \Omega^4_{BP, BP}(BP_*) \). Afterward, we will analyze \( d(\hat{x}) \), where \( d : \Omega^4_{BP, BP}(BP_*) \to \Omega^5_{BP, BP}(BP_*) \) denotes the differential in the cobar complex \( \Omega^*_{BP, BP}(BP_*) \). This analysis will provide us with the necessary information about \( d(\hat{x}) \), which represents \( d^\text{alg}_2(x) \in \text{Ext}_{P_*}^5(F_p, I/I^2) \).
Using Lemma 4.1 and the Leibniz rule, we have

\[
d(t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}}) = d(t_{4}^{p_{i}}) \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}} - t_{4}^{p_{i}} \otimes d(t_{3}^{p_{i}}) \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}}
\]
\[
+ t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes d(t_{1}^{p_{i}}) \otimes t_{2}^{p_{i}} - t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes d(t_{1}^{p_{i}})
\]
\[
\equiv R + pb_{4,i-1} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}} + L \mod I^2 \cdot BP_{*}BP^{\otimes 5}
\]
\[
\equiv R \mod I \cdot BP_{*}BP^{\otimes 5}
\]

where we denote

\[
R = (t_{3}^{p_{i}} \otimes t_{1}^{p_{i}+3} + t_{2}^{p_{i}} \otimes t_{1}^{p_{i}+2} + t_{1}^{p_{i}} \otimes t_{1}^{p_{i}+1}) \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}}
\]
\[
- t_{4}^{p_{i}} \otimes (t_{2}^{p_{i}} \otimes t_{1}^{p_{i}+2} + t_{1}^{p_{i}} \otimes t_{2}^{p_{i}+1}) \otimes t_{3}^{p_{i}} \otimes t_{1}^{p_{i}} + t_{1}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{1}^{p_{i}} \otimes t_{1}^{p_{i}+1} \otimes t_{1}^{p_{i}}
\]

and

\[
L = -t_{4}^{p_{i}} \otimes pb_{3,i-1} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}} + t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes pb_{2,i-1} \otimes t_{1}^{p_{i}} - t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes pb_{1,i-1}
\]

which is a sum of monomials in $I \cdot BP_{*}BP^{\otimes 5}$ with May degrees lower than

\[
M(pb_{4,j-1} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}}) = 7p + 10.
\]

Since $x \in Ext_{BP_{*}}^{4}(\mathbb{F}_{p}, BP_{*}/I)$ has May filtration leading term $t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}}$, we can choose a representative $\tilde{x}$ of $x$ in $\Omega_{BP_{*}}^{4}(BP_{*}) = BP_{*}BP^{\otimes 4}$ in the form of

\[
\tilde{x} = t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}} - \sum_{r} y_{r},
\]

such that:

(a) each $y_{r}$ is a monomial in $BP_{*}BP^{\otimes 4}$ and is not an element of $I \cdot BP_{*}BP^{\otimes 4}$,
(b) $M(y_{r}) < M(t_{4}^{p_{i}} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}}) = 7 + 5 + 3 + 1 = 16$,
(c) $\sum_{r} d(y_{r}) \equiv R \mod I^2 \cdot BP_{*}BP^{\otimes 5}$, ensuring that $d(\tilde{x}) \equiv 0 \mod I \cdot BP_{*}BP^{\otimes 5}$.

For each $r$, we express $d(y_{r})$ as a sum of monomials in $BP_{*}BP^{\otimes 5}$:

\[
d(y_{r}) = \sum_{u} z_{r,u}.
\]

Next, we define the sets $A_{r} := \{z_{r,u} | z_{r,u} \notin I \cdot BP_{*}BP^{\otimes 5}\}$ and $B_{r} := \{z_{r,u} | z_{r,u} \in I \cdot BP_{*}BP^{\otimes 5}, z_{r,u} \notin I^2 \cdot BP_{*}BP^{\otimes 5}\}$, which correspond to the (possibly empty) sets of summands. Using these sets, we then obtain:

\[
0 \equiv d(\tilde{x}) \equiv R - \sum_{r} \sum_{z_{r,u} \in A_{r}} z_{r,u} \mod I \cdot BP_{*}BP^{\otimes 5}
\]
\[
d(\tilde{x}) \equiv pb_{4,i-1} \otimes t_{3}^{p_{i}} \otimes t_{2}^{p_{i}} \otimes t_{1}^{p_{i}} + L - \sum_{r} \sum_{z_{r,u} \in B_{r}} z_{r,u} \mod I^2 \cdot BP_{*}BP^{\otimes 5}
\]
Moreover, these differentials are all nontrivial. We can take an example to show that $d''_2 (x) \in \text{Ext}^{5,5}_F (F_p, I/I^2)$.

The condition $M(y_r) < 16$ strongly restricts the form of $y_r$. To show that $M(z_{r,u}) < M(pb_{4,i-1} \otimes t''_1 \otimes t''_2 \otimes t''_{i+1}) = 7p + 10$ holds for all $z_{r,u} \in B_r$, we can conduct a tedious but straightforward check through all possible forms of $y_r$. Alternatively, we can summarize the idea as follows, considering three different cases:

(a) If $y_r = t''_4 \otimes A$ with $k \geq 1$, where $A$ is made up of $t_1$, $t_2$, and $t_3$ terms and $M(A) \leq 8$, then we have $M(z_{r,u}) \leq M(pb_{4,k-1} \otimes A) = 7p + 1 + M(A) \leq 7p + 9 < 7p + 10$.

(b) If $y_r = t''_4 \otimes A$, where $A$ is made up of $t_1$, $t_2$, and $t_3$ terms, and $M(A) \leq 8$, we note that $d(t''_4) = t''_3 \otimes t''_2 + t''_2 \otimes t''_1 - v_0 b_{3,0} - v_2 b_{2,1} - v_3 b_{1,2}$. Also, $M(b_{i,j}) = p(2i - 1) \leq 5p$ for $i \leq 3$. We can then observe that $M(z_{r,u}) < 7p + 10$.

(c) If $y_r$ is made up of $t_1$, $t_2$, and $t_3$ terms, we can also use similar ideas and check that $M(z_{r,u}) < 7p + 10$.

Thus, we conclude $d''_2 (x)$ has May filtration leading term $q_0 b_{4,i-1} \otimes t''_3 \otimes t''_2 \otimes t''_{i+1}$.

We can compute the following differentials similarly to Proposition 4.2.

**Proposition 4.3.** We have the following secondary algebraic Novikov differentials.

1. $d''_2 (t''_4 \otimes t''_3 \otimes t''_2 \otimes t''_{i+1}) = q_0 b_{4,i-1} \otimes t''_3 \otimes t''_2 \otimes t''_{i+1}$, for $i \geq 1$.
2. $d''_2 (t''_4 \otimes t''_3 \otimes t''_1 \otimes t''_{i+1}) = q_0 b_{4,i-1} \otimes t''_3 \otimes t''_1 \otimes t''_{i+1}$, for $i \geq 1$.
3. $d''_2 (t''_3 \otimes t''_2 \otimes t''_2 \otimes t''_{i+1}) = q_0 b_{3,i-1} \otimes t''_3 \otimes t''_2 \otimes t''_{i+1}$, for $i \geq 1$.

Here, the equations hold after modding out lower May filtration terms.

**Proof.** These results can be computed directly analogous to Proposition 4.2.

**Theorem 4.4.** There are nontrivial secondary Adams differentials given as follows:

1. $d''_{\text{Adams}} (h_{4,i} h_{3,i} g_i) = a_0 b_{4,i-1} h_{3,i} g_i$, for $i \geq 1$.
2. $d''_{\text{Adams}} (h_{4,i} h_{3,i+1} k_{i+2}) = a_0 b_{4,i-1} h_{3,i+1} k_{i+2}$, for $i \geq 1$.
3. $d''_{\text{Adams}} (h_{3,i} g_i h_{i+3}) = a_0 b_{4,i-1} h_{3,i+1} k_{i+3}$, for $i \geq 1$.
4. $d''_{\text{Adams}} (h_{3,i} h_{2,i+1} k_i) = a_0 b_{3,i-1} h_{2,i+1} k_i$, for $i \geq 1$.

**Proof.** These results can be directly deduced from Propositions 4.2 and 4.3. Moreover, these differentials are all nontrivial. We can take $a_0 b_{4,i-1} h_{3,i} g_i$ as an example to show $a_0 b_{4,i-1} h_{3,i} g_i \neq 0 \in \text{Ext}^{h_{4,i}}_{\text{Adams}}$. The other three cases are similar.
Note $a_0 b_{4,j-1} h_{3,j} g_i$ has May spectral sequence representative

$$a_0 b_{4,j-1} h_{3,j} h_{2,i} h_{1,i} \in E_1^{s,t,M}$$

Here, the inner degree is

$$t = 1 + q p^i((1 + p + p^2 + p^3) + (1 + p + p^2) + (1 + p) + 1)$$

where we denote $q = 2(p - 1)$. Let $x$ be an element in $E_1^{s,t,*}$. Inspection of degrees shows $x$ must be $a_0 h_{4,i} h_{3,j} h_{2,i} h_{1,i}$. Then $M(x) < M(a_0 b_{4,i-1} h_{3,j} h_{2,i} h_{1,i})$. Hence, $a_0 b_{4,j-1} h_{3,j} h_{2,i} h_{1,i}$ can not be the image of any May differential $d_r : E_r^{s,t,M + r} \to E_r^{s,t,M}$, $r \geq 1$. This completes the proof. \hfill \Box

It is worth pointing out that Zhong-Hong-Zhao [19] also computed two other nontrivial differentials on the fourth line.

**Theorem 4.5** (Zhong-Hong-Zhao [19]). On the fourth line $E_1^{5,*,*}$ of the Adams spectral sequence, there exist two nontrivial secondary Adams differentials given as follows:

1. $a_2^{Adams}(h_{3,i} g_i h_{2,j-1}) = a_0 b_{3,j-1} g_i h_{2,j-1}$ for $i \geq 2$.
2. $a_2^{Adams}(h_{3,i} k_{i+1} h_{2,j+2}) = a_0 b_{3,j+1} k_{i+1} h_{2,j+2}$ for $i \geq 1$.

Their result can be recovered by computing the following corresponding algebraic Novikov differentials.

**Proposition 4.6.** We have the following secondary algebraic Novikov differentials. Here, the equations hold after modding out lower May filtration terms.

1. $a_2^{alg}(t_3^i t_2^{p^i} t_1^{p^i} t_2^{p^{i-1}}) = q_0 b_{3,i-1} t_2^{p^i} t_1^{p^i} t_2^{p^{i-1}}$, for $i \geq 2$.
2. $a_2^{alg}(t_3^{p^{i+1}} t_2^{p^{i+1}} t_1^{p^{i+2}} t_2^{p^{i+2}}) = q_0 b_{3,j-1} t_2^{p^{i+1}} t_1^{p^{i+2}} t_2^{p^{i+2}}$, for $i \geq 1$.

**Proof.** These results can be computed directly analogous to Proposition 4.2. \hfill \Box

Our computations here are comparatively more straightforward than the original computations in [19] using matrix Massey products.

**5. Secondary Adams differentials on the first three lines**

In this section, we use the strategy explained in Section 4 to recover secondary Adams differentials on the first three lines.

The generators for the first two lines of the Adams spectral sequence were determined by Liulevicius in [7]. We summarize them in the following table.

| Generator | Representation in MSS | Inner Degree | Range of indices |
|-----------|-----------------------|--------------|-----------------|
| $a_0$     | $a_0$                 | 1            |                 |
| $h_i$     | $h_{1,i}$             | $qp^i$       | $i \geq 0$     |
| $a_1 h_0$ | $a_1 h_{1,0}$         | $2q + 1$     |                 |
| $a_0^2$   | $a_0^2$               | 2            |                 |
For odd primes, Aikawa [1] determined a basis for $\text{Ext}^{1,*}_{A_*}$ and $\text{Ext}^{2,*}_{A_*}$ using $\Lambda$-algebra. For $p \geq 5$, Wang [17] determined the May spectral sequence representatives of the generators. The result is summarized in the following table.

| Generator | MSS Representation | Inner Degree | Range of indices |
|-----------|--------------------|--------------|------------------|
| $h_i h_j$ | $h_{1,i} h_{1,j}$  | $q(p^i + p^j)$ | $k - 4 \geq j - 2 \geq i \geq 0$ |
| $a_i h_j$ | $a_{0,i} h_{1,j}$   | $q(p^i + p^j)$ | $j - 2 \geq i \geq 1$ |
| $a_i^2 h_j$ | $a_{0,i}^2 h_{1,j}$ | $q(p^i + p^j)$ | $i \geq 1$ |
| $a_0^2$ | $a_0^2$ | $3$ | |
| $b_i h_j$ | $b_{1,i} h_{1,j}$   | $q(p^i + p^j)$ | $i, j \geq 0, j \neq i + 2$ |
| $g_i a_0$ | $h_{2,i} a_{0,i}$   | $q(p^i + p^j)$ | $i \geq 1$ |
| $k_i h_j$ | $h_{2,i} h_{1,j}$   | $q(p^i + p^j)$ | $j \neq i + 2, i, i - 1, i \geq 0$ |
| $k_i a_0$ | $h_{2,i} a_{0,i}$   | $q(p^i + p^j)$ | $j \geq 0$ |
| $a_1 g_0$ | $a_{1,0} h_{1,0}$   | $q(p^i + p^j)$ | $i \geq 0$ |
| $h_{3,i} h_{1,i+2}$ | $h_{3,i} h_{1,i+2} h_{1,1}$ | $q(p^i + p^j)$ | $i \geq 0$ |
| $h_{3,i} k_{i+2}$ | $h_{3,i} h_{1,i+2} h_{1,1}$ | $q(p^i + p^j)$ | $i \geq 0$ |
| $a_1^2 h_0$ | $a_{1,0}^2 h_{1,0}$ | $3q + 2$ | |
| $b_{2,i} h_{1,i+1}$ | $b_{2,i} h_{1,i+1}$ | $q(p^i + p^j)$ | $i \geq 0$ |
| $b_{2,i} h_{1,i+2}$ | $b_{2,i} h_{1,i+2}$ | $q(p^i + p^j)$ | $i \geq 0$ |

Table 2. A $\mathbb{F}_p$-basis of $\text{Ext}^{1,*}_{A_*}$ and $\text{Ext}^{2,*}_{A_*}$

Table 3. A $\mathbb{F}_p$-basis of $\text{Ext}^{3,*}_{A_*}$
We can compute $d^2_{\text{Adams}}$ for the basis elements in Table 2 via computing $d^2_{\text{alg}}$ of their corresponding elements. For simplicity, we only list the nontrivial $d^2_{\text{alg}}$ differentials here.

**Proposition 5.1.** Let $p$ be an odd prime. Amongst the elements in the algebraic Novikov spectral sequence that corresponds to the first and second line basis listed in Table 2, all nontrivial $d^2_{\text{alg}}$'s are summarized as follows. Here, the equations hold after modding out lower May filtration terms.

1. $d^2_{\text{alg}}(t^p_j) = q_0b_{1,i-1}$, for $i > 0$.
2. $d^2_{\text{alg}}(pt^p_j) = q_0^2b_{1,i-1}$, $i \geq 1$.
3. $d^2_{\text{alg}}(t^p_j \otimes t^p_k) = q_0b_{2,i-1} \otimes t^p_k$, $i \geq 1$.
4. $d^2_{\text{alg}}(t_2 \otimes t_1) = -q_1b_{1,0} \otimes t_1$.
5. $d^2_{\text{alg}}(t_2 \otimes t^{p+1}_1) = q_0b_{2,i-1} \otimes t^{p+1}_1$, $i \geq 1$.
6. $d^2_{\text{alg}}(t^p_j \otimes t^p_k) = q_0b_{1,i-1} \otimes t^p_k - t^p_k \otimes q_0b_{1,j-1}$, $j - 2 \geq i \geq 1$.

**Proof.** Analogous to Proposition 4.2, all of the results are computed directly from the construction of the cobar complex. 

Then, we can recover the $d^2_{\text{Adams}}$ results on the first two lines directly from Proposition 5.1.

**Theorem 5.2** (Liulevicius[7], Shimada-Yamanosita [16], Miller-Ravenel-Wilson [12], Zhao-Wang [18]). Amongst the first and second line basis in Table 2, all nontrivial Adams $d^2$ differentials can be summarized as follows.

1. $d^2_{\text{Adams}}(h_i) = a_0b_{i-1}$, $i \geq 1$.
2. $d^2_{\text{Adams}}(a_0h_i) = a_0^2b_{i-1}$, $i \geq 1$.
3. $d^2_{\text{Adams}}(g_i) = a_0b_{2,i-1}h_i$, $i \geq 1$.
4. $d^2_{\text{Adams}}(g_0) = -a_1b_0h_0$.
5. $d^2_{\text{Adams}}(k_i) = a_0b_{2,i-1}h_{i+1}$, $i \geq 1$.
6. $d^2_{\text{Adams}}(h_ih_j) = a_0b_{1,i-1}h_j - h_ia_0b_{j-1}$, $j - 2 \geq i \geq 1$.

Similarly, we can compute $d^2_{\text{Adams}}$ for the third line basis via computing $d^2_{\text{alg}}$ of their corresponding elements. For simplicity, we only list the nontrivial differentials here.

**Proposition 5.3.** Let $p \geq 5$ be an odd prime. Amongst the elements in the algebraic Novikov spectral sequence that corresponds to the third line basis listed in Table 3, all nontrivial $d^2_{\text{alg}}$'s are summarized as follows. Here, the equations hold after modding out lower May filtration terms.

1. $d^2_{\text{alg}}(t^p_j \otimes t^p_k \otimes t^p_l) = q_0b_{1,i-1} \otimes t^p_l - t^p_l \otimes q_0b_{1,j-1} \otimes t^p_k + t^p_l \otimes t^p_k \otimes q_0b_{1,k-1}$, for $k - 4 \geq j - 2 \geq i \geq 1$.
2. $d^2_{\text{alg}}(q_0t^p_j \otimes t^p_k) = q_0^2b_{1,i-1} \otimes t^p_j - q_0^2t^p_j \otimes b_{1,j-1}$, for $j - 2 \geq i \geq 1$. 
(3) $d_{2}^{\text{alg}}(q_{0}^{2}t_{p}^{i}) = q_{0}^{3}b_{1,j-1}$, for $i \geq 1$.

(4) $d_{2}^{\text{alg}}(b_{0}t_{p}^{i}) = q_{0}^{3}b_{1,j-1}$, for $i \geq 0, j \geq 1, j \neq i + 2$.

(5) $d_{2}^{\text{alg}}(t_{2}^{i} \otimes t_{1}^{p} \otimes t_{1}^{j}) = q_{0}^{3}b_{2,i-1} \otimes t_{1}^{j}$, for $i \geq 1, j \geq 1, j 
eq i + 2, i, i - 1$.

(6) $d_{2}^{\text{alg}}(t_{2} \otimes t_{1} \otimes t_{1}^{j}) = -q_{1}b_{1,0} \otimes t_{1}^{j} + t_{2} \otimes t_{1} \otimes q_{0}b_{1,j-1}$, for $j > 0, j \neq 2$.

(7) $d_{2}^{\text{alg}}(q_{0}^{2}t_{p}^{i} \otimes t_{1}^{j}) = q_{0}^{3}b_{2,i-1} \otimes t_{1}^{j}$, for $i \geq 1$.

(8) $d_{2}^{\text{alg}}(t_{2}^{i} \otimes t_{1}^{j} + t_{1}^{j}) = q_{0}^{3}b_{2,i-1} \otimes t_{1}^{j}$, for $i \geq 1, j \geq 1, j \neq i + 2, i, i - 1$.

(9) $d_{2}^{\text{alg}}(q_{0}^{2}t_{p}^{i} \otimes t_{1}^{j}) = q_{0}^{3}b_{2,i-1} \otimes t_{1}^{j}$, for $i \geq 1$.

(10) $d_{2}^{\text{alg}}(t_{3}^{i} \otimes t_{2}^{j} \otimes t_{1}^{j}) = q_{0}^{3}b_{3,i-1} \otimes t_{2}^{j} \otimes t_{1}^{j}$, for $i \geq 1$.

(11) $d_{2}^{\text{alg}}(t_{3} \otimes t_{2} \otimes t_{1}) = -q_{1}b_{2,0} \otimes t_{1}$.

(12) $d_{2}^{\text{alg}}(t_{2} \otimes t_{2}^{i} \otimes t_{1}^{j}) = q_{0}^{3}b_{2,i-1} \otimes t_{2}^{i} \otimes t_{1}^{j}$, for $i \geq 1$.

(13) $d_{2}^{\text{alg}}(q_{1}t_{2} \otimes t_{1}) = q_{1}^{3}b_{1,0} \otimes t_{1}$.

(14) $d_{2}^{\text{alg}}(t_{3} \otimes t_{1}^{i} \otimes t_{1}^{j}) = q_{0}^{3}b_{3,i-1} \otimes t_{1}^{j}$, for $i \geq 1$.

(15) $d_{2}^{\text{alg}}(t_{3} \otimes t_{1}^{i} \otimes t_{1}) = -q_{1}b_{2,0} \otimes t_{1}^{i}$.

(16) $d_{2}^{\text{alg}}(t_{3} \otimes t_{1}^{i} \otimes t_{1}^{j}) = q_{0}^{3}b_{3,i-1} \otimes t_{2}^{i} \otimes t_{1}^{j}$, for $i \geq 1$.

Then, we can recover the following result directly from Proposition 5.3.

**Theorem 5.4** (Wang [17]). Let $p \geq 5$ be an odd prime. Amongst the third line basis in Table 3, all nontrivial Adams $d_{2}$ differentials can be summarized as follows.

(1) $d_{2}^{\text{Adams}}(h_{1}h_{1}h_{k}) = a_{0}b_{1,i-1}h_{1}h_{k} - a_{0}h_{1}b_{1,i-1}h_{k} + a_{0}h_{1}b_{1,k-1}$, for $k - 4 \geq j - 2 \geq i \geq 1$.

(2) $d_{2}^{\text{Adams}}(a_{0}h_{1}h_{j}) = a_{0}^{2}b_{1,i-1}h_{1} - a_{0}^{2}h_{1}b_{1,i-1}$, for $j - 2 \geq i \geq 1$.

(3) $d_{2}^{\text{Adams}}(a_{0}^{2}h_{j}) = a_{0}^{2}b_{1,j-1}$, for $j \geq 1$.

(4) $d_{2}^{\text{Adams}}(b_{1}h_{j}) = a_{0}b_{1,i-1}$, for $j \geq 0, j \geq 1, j \neq i + 2$.

(5) $d_{2}^{\text{Adams}}(g_{1}h_{j}) = a_{0}b_{1,i-1}h_{1}h_{j}$, for $j \geq 1, j \neq i + 2, i, i - 1$.

(6) $d_{2}^{\text{Adams}}(g_{0}h_{j}) = -a_{1}b_{0,j}h_{j} + a_{0}g_{0}b_{1,j-1}$, for $j > 0, j \neq 2$.

(7) $d_{2}^{\text{Adams}}(g_{1}a_{0}) = a_{0}^{2}b_{1,i-1}h_{i}$, for $i \geq 1$.

(8) $d_{2}^{\text{Adams}}(k_{1}h_{j}) = a_{0}b_{1,i-1}h_{1}h_{j}$, for $i \geq 1, j \neq i + 2, i, i - 1$.

(9) $d_{2}^{\text{Adams}}(k_{1}a_{0}) = a_{0}^{2}b_{2,i-1}h_{i+1}$, for $i \geq 1$.

(10) $d_{2}^{\text{Adams}}(h_{3,j}g_{i}) = a_{0}b_{3,j-1}g_{i}$, for $i \geq 1$.

(11) $d_{2}^{\text{Adams}}(h_{3,0}g_{0}) = -a_{1}b_{2,0}g_{0}$.

(12) $d_{2}^{\text{Adams}}(h_{2,i}g_{i+1}) = a_{0}b_{2,i-1}g_{i+1} - a_{0}h_{2,j}k_{i}$, for $i \geq 1$.

(13) $d_{2}^{\text{Adams}}(a_{1}g_{0}) = -a_{1}b_{0}h_{0}$.

(14) $d_{2}^{\text{Adams}}(h_{3,j}h_{i+2}) = a_{0}b_{3,j-1}h_{i+2}$, for $i \geq 1$.

(15) $d_{2}^{\text{Adams}}(h_{3,0}h_{2}h_{0}) = -a_{1}b_{2,0}h_{2}h_{0}$.

(16) $d_{2}^{\text{Adams}}(h_{3,j}k_{i+1}) = a_{0}b_{3,j-1}k_{i+1}$, for $i \geq 1$.  


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