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Preliminary investigation of the relationship between differential phase shift and path-integrated attenuation at the X band frequency in an Alpine environment

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Abstract. The RadAlp experiment aims at developing advanced methods for rainfall and snowfall estimation using weather radar remote sensing techniques in high mountain regions for improved water resource assessment and hydrological risk mitigation. A unique observation system has been deployed since 2016 in the Grenoble region of France. It is composed of an X-band radar operated by Météo-France on top of the Moucherotte mountain (1901 m above sea level; hereinafter MOUC radar). In the Grenoble valley (220 m above sea level; hereinafter a.s.l.), we operate a research X-band radar called XPORT and in situ sensors (weather station, rain gauge and disdrometer). In this paper we present a methodology for studying the relationship between the differential phase shift due to propagation in precipitation ($\Phi_{dp}$) and path-integrated attenuation (PIA) at X band. This relationship is critical for quantitative precipitation estimation (QPE) based on polarimetry due to severe attenuation effects in rain at the considered frequency. Furthermore, this relationship is still poorly documented in the melting layer (ML) due to the complexity of the hydrometeors’ distributions in terms of size, shape and density. The available observation system offers promising features to improve this understanding and to subsequently better process the radar observations in the ML. We use the mountain reference technique (MRT) for direct PIA estimations associated with the decrease in returns from mountain targets during precipitation events. The polarimetric PIA estimations are based on the regularization of the profiles of the total differential phase shift ($\Psi_{dp}$) from which the profiles of the specific differential phase shift on propagation ($K_{dp}$) are derived. This is followed by the application of relationships between the specific attenuation ($k$) and the specific differential phase shift. Such $k$–$K_{dp}$ relationships are estimated for rain by using drop size distribution (DSD) measurements available at ground level. Two sets of precipitation events are considered in this preliminary study, namely (i) nine convective cases with high rain rates which allow us to study the $\Phi_{dp}$–PIA relationship in rain, and (ii) a stratiform case with moderate rain rates, for which the melting layer (ML) rose up from about 1000 up to 2500 m a.s.l., where we were able to perform a horizontal scanning of the ML with the MOUC radar and a detailed analysis of the $\Phi_{dp}$–PIA relationship in the various layers of the ML. A common methodology was developed for the two configurations with some specific parameterizations. The various sources of error affecting the two PIA estimators are discussed, namely the stability of the dry weather mountain reference targets, radome attenuation, noise of the total differential phase shift profiles, contamination due to the differential phase shift on backscatter and relevance of the $k$–$K_{dp}$ relationship derived from DSD measurements, etc. In the end, the rain case study indicates that the relationship between MRT-derived PIA and polarimetry-derived PIA presents an overall coherence but quite a considerable dispersion (explained variance of 0.77). Interestingly, the nonlinear $k$–$K_{dp}$ relationship derived from independent DSD measurements yields almost unbiased PIA estimates. For the stratiform case, clear signatures of the MRT-derived PIA, the corresponding $\Phi_{dp}$ value and their ratio are evidenced within the ML. In particular, the av-
1 Introduction

Estimation of atmospheric precipitation (solid/liquid) is important in a high mountain region such as the Alps for the assessment and management of water and snow resources for drinking water, hydropower production, agriculture and tourism characterized by high seasonal variability. One of the most critical applications concerns the prediction of natural hazards associated with intense precipitation and melting of snowpacks, i.e., inundations, floods, flash floods and gravitational movements, which requires a high-resolution observation, namely spatial resolution $\leq 1$ km$^2$ and temporal resolution $\leq 1$ h. While this can hardly be achieved over extended areas with traditional in situ rain gauge networks, the use of radar remote sensing has a high potential that needs to be exploited but also a number of limitations that need to be surpassed. Quantitative precipitation estimation (QPE) with radar remote sensing in a complex terrain such as the Alps is made challenging by the topography and the space–time structure and the dynamics of precipitation systems. Radar coverage of the mountain regions brings the following dilemma. On the one hand, installing radar at the top of a mountain allows for a 360° panoramic view and therefore the ability to detect precipitation systems over a long range at the regional scale. This is particularly relevant for localized and heavy convective systems in warm seasons. But the precipitation is likely to undergo a significant change between detection and arrival at ground level, including a phase change when the 0°C isotherm is located at the level of or lower than the radar beam altitude. Such situations are likely to be frequent during cold periods, with a strong impact on QPE quality at ground level. On the other hand, installing a radar at the bottom of the valley provides high-resolution and quality data required for vulnerable and densely populated Alpine valleys, but the QPEs are limited in the latter due to beam blockage by surrounding mountains.

In Europe, MétéoSwiss has the longest-standing experience in operating radars in mountainous regions. The Swiss C-band radar network in the Alps (Joss and Lee, 1995; Germann et al., 2006) is one of the highest in the world and is coping with the associated altitude dilemma by using a large number of plan-position indicator (PPI) scans (including negative elevation ones) aimed at determining high-resolution vertical profiles of reflectivity. Sophisticated radar–rain gauge merging techniques and echo-tracking techniques, as well as numerical prediction models outputs (Sideris et al., 2014; Foresti et al., 2018) are implemented to better understand and quantify the complexity of precipitation distribution in such a rugged environment. More recently, Météo-France has chosen to complement the coverage of its operational radar network of Application Radar à la Météorologie Infra-Synoptique (ARAMIS) in the Alps by means of X-band polarimetric radars. A first set of three radars was installed in the southern Alps within the Risques Hydrométéorologiques en Territoires de Montagnes et Méditerranéens (RHYTMME) project in the period 2008–2013 at Montagne de Maurel (1770 m above sea level; hereinafter a.s.l.), Mont Colombis (1740 m a.s.l.) and Vars Mayt (2400 m a.s.l.; Westrelin et al., 2012). This effort was continued in 2014–2015 with the installation of an additional X-band radar system (hereinafter MOUC radar) on top of the Moucherotte mountain (1920 m) that dominates the valley of Grenoble, the biggest city in the French Alps with about 500 000 inhabitants. The choice of the X-band frequency is challenging due to its sensitivity to attenuation (e.g., Delrieu et al., 2000). In the past, the Institute of Environmental Geosciences (IGE) radar team has proposed the so-called mountain reference technique (MRT; Delrieu et al., 1997; Serrar et al., 2000; Bouilloud et al., 2009) to take advantage of this drawback for both correcting the gate-to-gate attenuation and performing a self-calibration of the radar. The idea was to estimate the path-integrated attenuations (PIAs) in some specific directions from the decrease in mountain returns during rainy periods. Such PIA estimates were then used as constraints for backward or forward attenuation correction algorithms (Marzoug and Amayenc, 1994) with optimization of an effective radar calibration error, given a drop size distribution (DSD) parameterization. The development of polarimetric radar techniques (e.g., Bringi and Chandrasekar, 2001; Ryzhkov et al., 2005) has allowed a scientific breakthrough for quantitative precipitation estimation (QPE) at X band by exploiting the relationship which exists between the specific differential phase shift on propagation ($K_{dp}$, in $^\circ$ km$^{-1}$) and the specific attenuation $k$ (dB km$^{-1}$). As with the MRT, the differential propagation phase $\Phi_{dp}(r_2) - \Phi_{dp}(r_1)$ over a given path ($r_1, r_2$) can be used to estimate $\text{PIA}(r_1, r_2)$, which can constrain a backward attenuation correction algorithm and allow a self-calibration of the radar and/or an adjustment of the DSD parameterization (Testud et al., 2000; Ryzhkov et al., 2014). Two major advantages of the polarimetric technique over the MRT can be formulated, namely (1) the availability of PIA constraints for any direction with significant precipitation and (2) the subsequent possibility of using a backward attenuation correction algorithm, which is known to be stable, while the forward formulation is inherently unstable. Accounting for their respective potential in different rain regimes (moderate to heavy), some combined algorithms...
making use of various polarimetric observables (reflectivity, differential reflectivity and specific differential phase shift on propagation) have also been proposed for the X-band frequency (e.g., Matrosov and Clark, 2002; Matrosov et al., 2005; Koffi et al., 2014). Although the polarimetric QPE methodology is now quite well established and validated for rainy precipitation (Matrosov et al., 2005; Anagnostou et al., 2004; Diss et al., 2009), Yu et al. (2018) have shown, in their first performance assessment of the RHyTMME radar network, the limitations associated with the use of polarimetric X-band radars in mountainous regions. They have pointed out (i) the need to better understand and quantify attenuation effects in the melting layer (ML), (ii) the importance of nonuniform beam filling (NUBF) effects at medium-to-long ranges in such a high-mountain context, and (iii) the stronger impact of radome attenuation at X band compared to S or C band. Yu et al. (2018) also had a first attempt at studying the relationship between the specific differential phase shift on propagation and the specific attenuation in the melting layer later by using the collocated measurements of two X-band radars, one situated well below and the other one situated well above the 0°C isotherm, and by considering the attenuation uniform within the ML.

Since 2016, we have had the opportunity to operate a research X-band polarimetric radar system (hereinafter XPORT radar) at IGE at the bottom of the Grenoble valley. This unique facility, consisting of two radar systems 11 km apart and operating on an altitudinal gradient of about 1700 m, should enable us to make progress on how to deal with the altitude dilemma and the potential/issues associated with the choice of the X-band operating frequency. Following a first article based on the RadAlp experiment about the characterization of the melting layer (Khanal et al., 2019), we concentrate hereinafter on the relationship between total differential phase shift ($\phi_{dp}$) derived from polarimetry and PIAs derived from the MRT. In Sect. 2, we present the observation system available and contrast rainy events considered in this study as follows: (i) a set of nine convective events with high rain rates, for which the melting layer was well above the detection domain of the XPORT radar, allows us to study the $\phi_{dp}$–PIA relationship in rain; and (ii) a stratiform case with moderate rain rates, for which the melting layer rose up from about 1000 to 2500 m a.s.l., allows us to perform a horizontal scanning of the ML with the MOUC radar and a preliminary analysis of the $\phi_{dp}$ – PIA relationship in the various layers of the ML. We present and illustrate, in Sect. 3, the methodology used for the PIA and $\phi_{dp}$ estimation. We also investigate the relationship between the specific differential phase shift on propagation ($K_{dp}$) and the specific attenuation ($k$) thanks to drop size distribution (DSD) measurements collected in the Grenoble valley during the two sets of events. The results concerning the $\phi_{dp}$–PIA relationship in rain and in the ML are presented and discussed in Sect. 4, while conclusions and perspectives are drawn in Sect. 5.

2 Observation system and datasets

2.1 Observation system

Grenoble is a y-shaped alluvial valley in the French Alps, with a mean altitude of about 220 m a.s.l. and surrounded by three mountain ranges, namely the Chartreuse (culminating at 2083 m a.s.l.) to the north, the Belledonne (2977 m) to the southeast and the Vercors (2307 m) to the west. Figure 1 shows the topography of the area and the positions of the Météo-France radar system on top of the Moucherotte mountain and the IGE experimental site at the bottom of the valley.

Among other devices, the IGE experimental site includes the following: (i) the IGE XPORT research radar (Koffi et al., 2014; see Table 1 for the list of its main parameters); (ii) one Micro Rain Radar (MRR; not used in the current study); (iii) one meteorological station including pressure, temperature, humidity, wind probes and several rain gauges; and (iv) one Parsivel$^2$ disdrometer. The characteristics of the MOUC radar are listed in Table 1. XPORT radar was built in the laboratory in the 2000s. It was operated during more than 10 years in western Africa within the African Monsoon Multidisciplinary Analysis (AMMA) and Megha-Tropiques calibration/validation campaigns. Since its return to France in 2016, a maintenance and updating program has been underway to improve its functionalities, notably with respect to the real-time data processing and the antenna control program. One noticeable feature of XPORT radar is the range bin size of 34.2 m (which actually corresponds to an oversampling since, for a pulse width of 1 μs, the theoretical bin size is 150 m), which is an interesting figure for the close-range and volumetric measurements considered in this study. Note that while the MOUC radar is operated 24 h per day and its data are integrated in the Météo-France mosaic radar products, the XPORT radar is operated with alerts only on significant precipitation events.

2.2 Dataset

Table 2 shows the main characteristics of the nine convective events considered for the study of the $\phi_{dp}$–PIA relationship in rain, by using the XPORT radar data. A stratiform event, which occurred on 3–4 January 2018, is also considered for a preliminary study of the $\phi_{dp}$–PIA relationship in the ML, with both the MOUC and the XPORT radar data. Figure 2 presents a time series of one of the most intense convective event (21 July 2017) and the stratiform event. In both cases, the total rain amount observed at the IGE site was about 35 mm, but in 3 h, with two peak rain rates of about 40 mm h$^{-1}$ for the 21 July 2017 convective event, while the 3–4 January 2018 stratiform event lasted more than 12 h with an average rain rate of about 3 mm h$^{-1}$. The two events also differ in their vertical structure. The bottom graphs of Fig. 2 display the time series of the altitudes of the tops, peaks...

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Figure 1. The topographical map of Grenoble is shown along with the positions of two radar systems. A vertical cross section along the line joining the two radar sites is shown in the insert at the bottom right of the figure.

Table 1. Characteristics of the XPORT and MOUC radar systems.

| Characteristics                  | MOUC radar              | XPORT radar             |
|----------------------------------|-------------------------|-------------------------|
| Longitude (decimal degrees)      | 5.639237                | 5.762327                |
| Latitude (decimal degrees)       | 45.147736               | 45.194150               |
| Altitude (m a.s.l.)              | Ground: 1901            | Ground: 213             |
| Antenna feed horn                | Antenna feed horn: 1917 | Antenna feed horn: 228  |
| Frequency (GHz)                  | 9.420                   | 9.400                   |
| Antenna diameter (m)             | 1.8                     | 1.8                     |
| 3 dB beamwidth (°)               | 1.28                    | 1.37                    |
| Antenna gain (dB)                | 42                      | 42                      |
| Radome                           | yes                     | no                      |
| Peak power                       | 30 kW on each polarization | 50 kW on each polarization |
| Pulse length (µs)                | 2                       | 1                       |
| Radial bin size (m)              | 240                     | 34.2                    |
| Receiver dynamic range (dBm)     | −114                    | −112                    |
| Minimum detectable signal (dBm)  | 0/0.6/1.2/2/3/4/8/14°   | 3.5/7.5/15/25/45°      |
| Volume scanning protocol (PPIs with elevation angles in °) | 5                       | ~ 7                     |
| Measured parameters              | $Z_h$, $Z_v$, $Z_{dr}$, $\rho_{hv}$, $\phi_{dp}$, $v_r$ | |

and bottoms of the horizontal reflectivity ($Z_h$) and copolar correlation coefficient ($\rho_{hv}$) signatures of the ML, obtained with the automatic detection algorithm described in Khanal et al. (2019). The quasi-vertical profiles (QVP; Ryzhkov et al., 2016) derived from the XPORT 25° PPIs are considered in the ML detection. For the convective case, the ML extends from 3000 to 4000 m a.s.l. and more, i.e., well above the altitudes of the two radars. Table 2 indicates that this is also the case for the other convective events – at least for the XPORT radar. For the stratiform event, the ML extends between 800 and 1500 m a.s.l. during the first part of the event (between 3 January, 20:00 UTC, and 4 January, 01:30 UTC) and then rises in about 2 h to stabilize at an altitude range of about
Table 2. Some characteristics of the nine convective events considered in the study of the $\phi_{dp}$–PIA relationship in rain. The ML detection was performed with the 25° elevation angle measurements of the XPORT radar using the algorithm described in Khanal et al. (2019). The total rain amount and the maximum rain rate are recorded at the weather station available at the XPORT radar site at IGE. The maximum PIA is derived from the MRT by considering the 7.5° elevation data of the XPORT radar.

| Date       | Beginning (UTC) | End (UTC) | Minimum altitude of ML bottom (m a.s.l.) | Total rain amount (mm) | Maximum rain rate in 10 min (mm h$^{-1}$) | Maximum MRT PIA (dB) |
|------------|-----------------|-----------|-------------------------------------------|------------------------|------------------------------------------|----------------------|
| 12 May 2017| 12:00           | 16:00     | 2000                                      | 9.2                    | 8.4                                      | 14.2                 |
| 21 July 2017| 15:30           | 19:30     | 3000                                      | 35.2                   | 42.0                                     | 30.7                 |
| 8 August 2017| 08:30           | 14:30     | 3700                                      | 27.9                   | 48.0                                     | 30.1                 |
| 31 August 2017| 07:00           | 11:30     | 3200                                      | 19.9                   | 15.5                                     | 7.6                  |
| 22 May 2018 | 16:00           | 23:00     | 2000                                      | 16.9                   | 8.4                                      | 10.2                 |
| 27 May 2018 | 14:00           | 17:00     | 2700                                      | 6.9                    | 9.9                                      | 6.0                  |
| 28 May 2017 | 13:00           | 23:00     | 2500                                      | 9.8                    | 9.0                                      | 7.7                  |
| 20 July 2018| 17:00           | 22:00     | 2700                                      | 12.1                   | 15.6                                     | 19.3                 |
| 9 August 2018| 07:30           | 15:00     | 3000                                      | 24.8                   | 8.4                                      | 19.2                 |

2200–2800 m a.s.l. after 04:00 UTC, passing progressively at the level of the MOUC radar in the meantime.

As an additional illustration of the dataset, Fig. 3 gives two examples of XPORT PPIs at 7.5° elevation angle for moderate (left) and intense (right) rain during the 21 July 2017 event. As a clear feature, one can see that, for this elevation angle, the radar beam is fully blocked by the Chartreuse mountain range in the northern sector. Also visible in the northeastern sector and, to a lesser extent, in the southwestern sector are partial beam blockages associated with tall trees in the vicinity of the XPORT radar on the Grenoble campus. This figure is also intended to draw the attention of the reader to the decrease in the Chamrousse and Moucherotte mountains returns (within red circles) during the intense rain time step, compared to their values in moderate rain, as a first illustration of the MRT principle.

### 3 Methodology

Our aim is to study the relationship between two radar observables of propagation effects at X band, namely path-integrated attenuation and differential propagation phase due to precipitation occurring along the radar path. We describe, in the following two subsections, the estimation methods that were implemented. In Sect. 3.3, we complement the methodology description with the presentation of DSD-derived $k$–$K_{dp}$ relationships.

#### 3.1 Path-integrated attenuation estimation

Let us express the PIAs (in dB) at a given range $r$ (km) as follows:

$$\text{PIA}(r) = \text{PIA}(r_0) + 2 \int_{r_0}^{r} k(s) \, ds,$$

(1)

where $k(s)$ (dB km$^{-1}$) is the specific attenuation due to rain at range $s$ (km). $r_0$ is the range where the measurements start to become exploitable, i.e., the range where measurements are free of ground clutter associated with side lobe effects. The term PIA($r_0$) represents the so-called on-site attenuation resulting from radome attenuation and range attenuation at range closer than $r_0$. Note that PIAs can be obtained from Eq. (1) for both the horizontal and the vertical polarizations. In the present article, we will restrict ourselves to the horizontal polarization, as the study of differential attenuation is a possible topic for a future study. Delrieu et al. (1999) have proposed an assessment of the quality of PIA estimates from mountain returns by implementing a receiving antenna in the Belledonne mountain range in conjunction with an X-band radar operated on the Grenoble campus. They found good agreement between the two PIA estimates for PIAs exceeding the natural variability of the mountain reference target during dry weather. They recommended using strong mountain returns (greater than, e.g., 50 dBZ during dry weather) so as to minimize the impact of precipitation falling over the reference target itself. They also point out that this approach is not able to separate the effects of on-site and range attenuation. They verified, however, by implementing the receiving antenna close to the radar (at a range of about 200 m), that the on-site attenuation was negligible for a radomeless radar, which is the case for the XPORT radar but not for the MOUC radar. Another interesting feature of the MRT PIA estimator is its independence with respect to eventual radar calibration errors.
Figure 2. Description of two rain events considered in the present study: (a, c) convective case of 21 July 2017 and (b, d) stratiform case of 4 January 2018. (a, b) Rain rate and cumulative rainfall time series observed at the IGE site and (c, d) results of the ML detection algorithm based on XPORT 25° PPI data. The horizontal red line indicates the altitude of the MOUC radar (see text for details).

Figure 3. Examples of XPORT 7.5° PPIs of raw reflectivity (not corrected for attenuation) taken for two time steps during the 21 July 2017 convective event. The crosses indicate the location of the two radars and the black/white (5/10 km) range markers correspond to the XPORT and the MOUC radar, respectively. The red circles highlight the mountain returns associated with the Chamrousse (southeast) and the Moucherotte (southwest) mountains in between the 10–15 km range.
In the current study, we used the following procedure to determine the mountain reference targets for the XPORT radar.

A large series of raw reflectivity data, observed during widespread rainfall with no ML contamination, was accumulated and averaged in order to characterize the detection domain of the XPORT radar at the 7.5° elevation angle. This allowed us to determine the mountain returns, the full beam blockages due to mountains, the partial beam blockages due to tall trees and spurious detections due to side lobes in the vicinity of the radar. A manual selection of the mountain reference targets was then performed based on the map of the apparent reflectivity above 45 dBZ. The targets, made up of mountain returns from successive radials (up to 9) with a limited range extent (less than 2.0 km), are described in Table 3. Based on the radar equation and the receiver characteristics, care was taken to discard targets eventually subject to saturation at close range. The selected targets are located at a mean range between 4.1 and 17.1 km and have sizes between 0.06 and 0.94 km². For each rain event, dry weather data before and/or after the event were used to characterize the mean target reflectivity and its time variability. Note that the mean reflectivity for each target and each time step was computed as the average of the dBZ values of each radial gate comprising the target. This is justified by the fact that we aim at estimating PIA in dB. Table 3 lists the mean, standard deviation and 10% and 90% quantiles of the time series of the dry weather apparent reflectivity of the reference targets for the first and last event of the considered series. One can notice the good stability of the mean reflectivity values between the two events, which is an indication of both the radar calibration stability during the period and the moderate impact of the mountain surface conditions already evidenced in previous studies (Delrieu et al., 1999; Serrar et al., 2000). The standard deviations of the reflectivity time series range between 0.2 and 0.9 dBZ, and the mean 10%–90% interquantile range is equal to 1.03 dBZ.

Due to limited data availability, a simpler approach was implemented for the selection of the MOUC mountain reference targets. Here again, the raw reflectivity data were accumulated and averaged, but only over the period of 3 January 2018, 19:00–23:55 UTC, preceding the rise of the ML at the level of the MOUC radar. It was snowing during this period at the MOUC radar site. So, we are implicitly making the assumption of negligible attenuation during snowfall (supported by the literature; e.g., Matrosov et al., 2009) in the considered case study. Table 4 displays the geometrical characteristics of the targets, as well as the mean, standard deviation and 10% and 90% quantiles of their apparent reflectivity time series. Targets are located at greater distances than those of the XPORT radar, i.e., between 19.9 and 44.9 km. In spite of having larger sizes (between 0.7 and 4.0 km²), this range effect probably explains why their standard deviations are higher, namely between 0.75 and 1.44 dBZ. The 10%–90% interquantile ranges are subsequently higher as well, with a mean value of 2.6 dBZ.

The top graphs of Fig. 4 give two examples of apparent reflectivity profiles for a radial of a given target during the 21 July 2017 rain event. The example on the left-hand side corresponds to a moderate PIA (5.4 dB, when considering all the gates of the radials comprising the target) and the right-hand side example corresponds to one of the highest PIA value observed (27.6 dB) in our dataset. We tried to limit, as far as possible, the radial extent of targets (less than 2000 m) and/or multipeaks targets, such as the one shown in the left-hand side example, in order to limit positive bias on MRT PIA estimates. The top graphs of Fig. 5 give two examples of apparent reflectivity time series during the events of 21 July 2017 and 20 July 2018, together with the mean, 10% and 90% quantiles of the dry weather apparent reflectivity. For both cases, the XPORT data acquisition started a bit after the actual beginning of the storm. Therefore, the dry weather reference values were estimated with data collected after the event, i.e., between 19:00 and 22:00 UTC for the 21 July 2017 event and between 00:00 and 06:00 UTC on the day after for the 20 July 2018 event. For these convective events, one can note the erratic nature of the apparent reflectivity time series at the XPORT radar acquisition period used at that time (about 7 min). The MRT PIA estimates are simply calculated as the difference between the mean values of the target apparent reflectivity during dry weather and at each time step of the rain event (blue lines in the bottom graphs of Fig. 5).

3.2 Differential propagation phase estimation

Let us express the total differential phase shift between copolar (HH and VV) received signals as follows:

$$\psi_{dp}(r) = 2 \int_{r_0}^{r} K_{dp}(s) \, ds + \delta_{hv}(r),$$  \hspace{1cm} (2)

where $$K_{dp}(s)$$ is the specific differential phase shift on propagation ($$\circ$$ km$$^{-1}$$) related to precipitation at any range $$s$$ between $$r_0$$ and $$r$$, and $$\delta_{hv}(r)$$ is the differential phase shift on backscatter ($$\circ$$) at range $$r$$.

The quantity of interest, i.e., the differential propagation phase associated with precipitation along the path, is denoted as follows:

$$\phi_{dp}(r) = 2 \int_{r_0}^{r} K_{dp}(s) \, ds = \psi_{dp}(r) - \delta_{hv}(r).$$  \hspace{1cm} (3)

As with the on-site attenuation for the MRT, we have a problem here with the possible influence of the differential phase shift on backscatter $$\delta_{hv}(r)$$ that may introduce a positive bias on the estimation of the differential phase shift associated with precipitation along the path. In the literature (e.g.,

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Table 3. Geometrical characteristics and apparent reflectivity statistics for the 16 mountain targets selected for the XPORT radar at an elevation angle of 7.5°. The mean, standard deviation, and 10% and 90% quantiles of the apparent reflectivity time series are given for the first and last convective events in the considered period (see Table 2).

| Target | Mean azimuth (°) | Mean range (km) | Number of gates | Size (km²) | Mean reflectivity (dBZ) | Standard deviation (dBZ) | 10% quantile (dBZ) | 90% quantile (dBZ) | 12 May 2017 | 9 August 2018 |
|--------|-----------------|----------------|----------------|-----------|------------------------|-------------------------|-------------------|-------------------|-------------|--------------|
| 1      | 2.9             | 4.1            | 51             | 0.06      | 48.39                  | 0.21                    | 48.17             | 48.73             | 48.29       | 48.03        |
| 2      | 13.2            | 4.8            | 130            | 0.18      | 52.19                  | 0.80                    | 51.55             | 53.17             | 52.22       | 51.62        |
| 3      | 17.5            | 5.7            | 163            | 0.27      | 51.90                  | 0.29                    | 51.61             | 52.25             | 52.42       | 51.91        |
| 4      | 24.0            | 8.6            | 133            | 0.33      | 51.98                  | 0.51                    | 51.44             | 52.80             | 51.87       | 51.41        |
| 5      | 29.0            | 14.6           | 71             | 0.30      | 49.44                  | 0.55                    | 48.91             | 50.01             | 50.31       | 49.63        |
| 6      | 89.5            | 17.1           | 160            | 0.79      | 53.20                  | 0.38                    | 52.81             | 53.59             | 52.78       | 52.34        |
| 7      | 95.3            | 14.5           | 95             | 0.40      | 54.12                  | 0.23                    | 53.91             | 54.30             | 53.96       | 53.72        |
| 8      | 98.4            | 13.2           | 120            | 0.45      | 51.02                  | 0.50                    | 50.59             | 51.67             | 52.13       | 51.69        |
| 9      | 101.2           | 13.1           | 156            | 0.58      | 48.95                  | 0.23                    | 48.71             | 49.18             | 49.50       | 49.37        |
| 10     | 119.7           | 12.1           | 92             | 0.32      | 49.36                  | 0.21                    | 49.11             | 49.59             | 50.23       | 50.07        |
| 11     | 124.8           | 11.8           | 242            | 0.82      | 51.04                  | 0.53                    | 50.50             | 51.84             | 52.02       | 51.63        |
| 12     | 130.1           | 11.9           | 240            | 0.82      | 51.43                  | 0.90                    | 50.48             | 52.63             | 54.63       | 54.08        |
| 13     | 135.1           | 12.0           | 271            | 0.94      | 52.97                  | 0.67                    | 52.20             | 53.59             | 53.24       | 53.78        |
| 14     | 228.8           | 11.4           | 221            | 0.73      | 52.63                  | 0.59                    | 51.97             | 53.46             | 53.79       | 53.37        |
| 15     | 228.8           | 10.7           | 187            | 0.58      | 53.62                  | 0.41                    | 53.11             | 54.02             | 52.96       | 52.53        |
| 16     | 248.8           | 10.5           | 162            | 0.49      | 53.56                  | 0.41                    | 53.11             | 54.02             | 52.96       | 52.53        |

Table 4. Geometrical characteristics and apparent reflectivity statistics for the 13 mountain targets selected for the MOUC radar at an elevation angle of 0°. The mean, standard deviation, and 10% and 90% quantiles of the apparent reflectivity time series are computed over the period of 3 January 2018 from 19:00 to 23:55 UTC, preceding the rising of the ML at the level of the MOUC radar.

| Target | Mean azimuth (°) | Mean range (km) | Number of gates | Size (km²) | Mean reflectivity (dBZ) | Standard deviation (dBZ) | 10% quantile (dBZ) | 90% quantile (dBZ) |
|--------|-----------------|----------------|----------------|-----------|------------------------|-------------------------|-------------------|-------------------|
| 1      | 40.0            | 29.52          | 25             | 1.55      | 49.97                  | 1.1                     | 48.70             | 51.26             |
| 2      | 43.7            | 26.28          | 13             | 0.72      | 49.90                  | 1.28                    | 48.18             | 51.36             |
| 3      | 78.0            | 27.12          | 24             | 1.36      | 48.18                  | 1.44                    | 46.77             | 50.12             |
| 4      | 84.2            | 23.64          | 28             | 1.39      | 49.56                  | 0.95                    | 48.37             | 50.90             |
| 5      | 89.5            | 23.04          | 82             | 3.96      | 49.09                  | 0.62                    | 48.39             | 49.80             |
| 6      | 96.0            | 21.36          | 78             | 3.49      | 49.37                  | 0.75                    | 48.32             | 50.34             |
| 7      | 101.7           | 19.92          | 52             | 2.17      | 49.31                  | 1.01                    | 47.83             | 50.37             |
| 8      | 107.2           | 22.44          | 33             | 1.55      | 51.94                  | 1.11                    | 50.52             | 53.22             |
| 9      | 117.0           | 25.32          | 38             | 2.02      | 51.50                  | 1.03                    | 50.17             | 52.74             |
| 10     | 121.2           | 23.52          | 41             | 2.02      | 48.65                  | 1.18                    | 47.28             | 50.18             |
| 11     | 128.5           | 28.44          | 43             | 2.56      | 49.38                  | 0.98                    | 48.21             | 50.59             |
| 12     | 132.5           | 27.00          | 25             | 1.41      | 50.33                  | 1.24                    | 48.71             | 51.81             |
| 13     | 160.2           | 44.88          | 37             | 3.48      | 49.91                  | 1.00                    | 48.69             | 51.11             |

Otto and Russchenberg, 2011; Schneebeli and Berne, 2012) we find power law relationships between δ_{hv} and Z_{dr} at X band in rain, giving values for the differential phase shift on backscatter in the ranges of [0.6–1.0°] and [2.1–3.5°] for the differential reflectivity of 1 and 2 dB, respectively. Scattering simulations based on disdrometer data (Trömel et al., 2013) indicate that quite a large scatter may exist with respect to such power law models, and there is an important influence from the considered hydrometeor temperature. From simulations based on radar data at various frequencies, the same authors quantify δ_{hv}(r) values as high as 4° in the ML at X band and mention that strong δ_{hv}(r) values may be associated with both large dry hailstones and wet hailstones, especially at X band. Let us note that no hail was reported for the convective cases considered in the present study. Keeping the related orders of magnitude in mind, and the fact that significant δ_{hv} effects are associated with bumps in the ψ_{dp}
profiles, hereafter we will carefully discuss the possibility of assuming $\delta_{hv}$ to be negligible, or not, with respect to $\phi_{dp}$.

In this study, the following method was implemented for the processing of the $\psi_{dp}$ profiles and the subsequent estimation of $\phi_{dp}$ values near the mountain targets for the XPORT radar (rain case based on convective events).

We first determined so-called rainy range gates along the path by using the $\rho_{hv}$ profiles. The raw $\psi_{dp}(r)$ values, for which $\rho_{hv}(r)$ was less than 0.95 (empirical threshold with limited impact in the [0.95–0.97] range), were set to missing values. In addition, we defined the beginning of the rainy range by determining the first series of 10 successive gates (again, an empirical choice corresponding to a range extent of 342 m) overpassing this threshold. The $r_0$ value was set to the minimum range value of this series. Similarly, we defined the end of the rainy range by determining the last series of 10 successive range gates overpassing this threshold close to the mountain target. A maximum rainy range, denoted as $r_M$, was defined as the maximum range value of this series. It is noteworthy to mention that rain likely occurs in the ranges less than $r_0$ and greater than $r_M$ and in the intermediate ranges for which the $\psi_{dp}(r)$ values were set to missing values. It is, however, critical to discard such gates that may be prone to clutter due to side lobes close to the radar or mountain returns close to the mountain target. Although the intermediate missing values will not impact the $\phi_{dp}$ estimation, we have to mention that both the initial and final missing values may result in a negative bias on the PIA estimation based on $\phi_{dp}(r_M)$.

In the current version of the procedure, every single radial was processed separately. First, an unfolding was applied by adding 360° to negative $\psi_{dp}(r)$ values. The system differen-
tial phase shift was estimated as the median of the $\psi_{dp}(r)$ values corresponding to the beginning of the rainy range. This value was subtracted from the raw $\psi_{dp}(r)$ profiles, and eventual negative values were set to 0. Regarding the $\psi_{dp}$ measurement noise processing, we have implemented and improved a regularization procedure initially proposed by Yu and Gaussiat (2018). This procedure consists of defining an upper envelope curve, starting from $r_0$, and a lower envelope curve, starting from $r_M$, by considering a maximum jump, denoted as $\text{diffmax}$, authorized between two successive gates. The calculation was performed for a series of $\text{diffmax}$ values in the range of 0.5–10°. The regularized $\psi_{dp}$ profiles (increasing monotonous curves) were estimated by taking the average of the upper and lower envelope curves. Note that the values for the missing gates between $r_0$ and $r_M$ were simply interpolated with the adjacent values of the regularized profile. A mean absolute difference (MAD) criterion between the raw and regularized profiles over a series of 30 gates with nonmissing values near the mountain target (an empirical choice corresponding to a range extent of about 1 km) was used to determine the optimal $\text{diffmax}$ value and the associated profile. The optimal profile was finally selected if the MAD criterion was less than 50%, otherwise we considered the polarimetry-derived PIAs to be missing for the considered radial. Finally, the $\phi_{dp}(r_M)$ value for the target was estimated as a weighted average of the $\phi_{dp}(r_M)$ values of all the nonmissing radials composing the target, with the weights being the number of reference gates of each radial. The bottom graphs of Fig. 4 present the raw and regularized profiles, and the envelope curves, for the examples already commented on above. For the right-side hand example corresponding to one of the strongest PIAs (27.6 dB) observed, one can note that the noise of the raw $\psi_{dp}$ profile is low, especially in the range with the highest gradients between 7 and 13 km. There is no apparent bump on the raw profile that could signal a $\delta_{hv}$ contamination so that one might be tempted to consider the regularized profile as a good estimator of the $\phi_{dp}$ profile in that case. The left-hand side example, corresponding to a moderate MRT-derived PIA of 5.4 dB, is more complex. As already noted, the mountain target itself is noisy with significant mountain return contamination before range $r_M$, as evidenced by the $\rho_{hv}$ profile. In addition, one can note a nonmonotonic behavior of the raw $\psi_{dp}$ profile, with a plateau of about 17.5° for ranges greater than 4 km, follow-
4 January 2018, target 7

(a) MOUC radar time series

(b) \( PIA(r_M) \) [dB]

(c) \( \phi_{dp}(r_M) \) [°]

(d) XPORT radar: melting layer identification

Figure 6. Time series of (a) the apparent reflectivity values of a given mountain reference target together with the dry weather reference (solid red horizontal lines for the mean) and the 10% and 90% quantiles (broken red lines), (b) the resulting PIA estimates (dB), and (c) the corresponding \( \phi_{dp}(r_M) \) values (°) for the 0° PPI of the MOUC radar during the 3–4 January 2018 stratiform rain event. (d) displays the results of the ML detection algorithm performed with the XPORT 25° PPI data (see text for details).

As representative examples, Fig. 7 illustrates the range profiles taken by the MOUC radar during the snowfall (left) and the ML (right) periods. As expected, the \( \rho_{hv} \) profiles are very different in the two cases, with \( \rho_{hv} \) values close to 1 in snow, indicating precipitation homogeneity while \( \rho_{hv} \) presents a high variability in the ML. During the ML period, we therefore had to adapt the \( \rho_{hv} \) threshold used to detect gates with precipitation. Based on the \( \rho_{hv} \) peak statistics presented by Khanal et al. (2019), we have chosen a value of 0.8. As it can be seen in Fig. 7, such a threshold may prevent the detection of the mountain reference return itself. Subsequently, we had to adapt the determination of ranges \( r_M \) and \( r_M \) with respect to the XPORT radar case, firstly, by considering two successive gates corresponding to a range extent of 480 m (instead of 10 gates corresponding to 342 m) and, secondly, by making sure that the calculated \( r_M \) value was less than the range of the first mountain reference gate. Regarding the regularization of the \( \psi_{dp} \) profiles (bottom graphs of Fig. 7), it was found that the raw profiles were noisier compared to the XPORT case study. Well-structured bumps were not evidenced in the ML profiles, maybe as a result of the lower-range resolution of the MOUC radar, and the regularization procedure was found to work satisfactorily. It remains, however, difficult to assume that there is no \( \delta_{hv} \) contamination during the ML period.

Coming back to Fig. 6, one can note the mean value of \( \phi_{dp}(r_M) \) to be equal to 11.2° during the snowfall period, resulting in a specific differential phase shift on a propagation of 0.28° km\(^{-1}\) if the differential phase shift on backscatter is neglected. Such values indicate a significant heterogeneity in the horizontal and vertical dimensions of the snow/ice hydrometeors. During the rainy period between 04:10 and 06:00 UTC, there is good coherence between the specific PIA estimates and the \( \phi_{dp}(r_M) \) estimates for the 0° PPI of the MOUC radar during the stratiform event of 3–4 January 2018. The time period considered in the figure ranges from 00:00 to 06:00 UTC on 4 January 2018 in order to focus on the rising of the ML between 02:00 and 04:00 UTC. The target is located at a distance of 19.9 km from the radar. The bottom graph of Fig. 6 displays the results of the ML detection algorithm (Khanal et al., 2019) in terms of the altitudes of the top, peak and bottom of the \( Z_h \) (blue) and the \( \rho_{hv} \) (orange) ML signatures. The altitude of the \( Z_h \) top inflexion point is assumed to correspond to the 0°C isotherm altitude, while the \( \rho_{hv} \) bottom inflexion point corresponds well with the bottom of the ML according to Khanal et al. (2019). We therefore define the ML width as the altitude difference between \( Z_h \) top and \( \rho_{hv} \) bottom. Before 02:00 UTC, the ML is well below the altitude of the MOUC radar. MOUC radar measurements at the 0° elevation angle are therefore made in snow/ice precipitation during this period. Based on the ML detection results, the passage of the ML at the altitude of the MOUC radar begins at about 02:20 UTC and ends at 04:10 UTC. After this time, MOUC radar measurements are therefore made in rainfall.

Basically, the same methodology was implemented for the MOUC radar case study, with some alterations to be described hereafter. Figure 6 provides the time series of the apparent reflectivity of a given mountain target, the resulting increase in the raw profile (with moderate noise) up to 22° at a 4 km range. One might assume that there is a \( \delta_{hv} \) contamination in that case. Interestingly, the regularization procedure is shown to provide a good filtering of the bump, and here again we are tempted to consider the regularized profile as a good estimator of the \( \phi_{dp} \) profile. The middle graphs in Fig. 5 display the time series of the \( \phi_{dp}(r_M) \) values associated with the apparent reflectivity of the mountain returns discussed above. One can note good consistency in the two time series for the highest peaks, while discrepancies can be evidenced for the moderate and small values.

As representative examples, Fig. 7 illustrates the range profiles taken by the MOUC radar during the snowfall (left) and the ML (right) periods. As expected, the \( \rho_{hv} \) profiles are very different in the two cases, with \( \rho_{hv} \) values close to 1 in snow, indicating precipitation homogeneity while \( \rho_{hv} \) presents a high variability in the ML. During the ML period, we therefore had to adapt the \( \rho_{hv} \) threshold used to detect gates with precipitation. Based on the \( \rho_{hv} \) peak statistics presented by Khanal et al. (2019), we have chosen a value of 0.8. As it can be seen in Fig. 7, such a threshold may prevent the detection of the mountain reference return itself. Subsequently, we had to adapt the determination of ranges \( r_M \) and \( r_M \) with respect to the XPORT radar case, firstly, by considering two successive gates corresponding to a range extent of 480 m (instead of 10 gates corresponding to 342 m) and, secondly, by making sure that the calculated \( r_M \) value was less than the range of the first mountain reference gate. Regarding the regularization of the \( \psi_{dp} \) profiles (bottom graphs of Fig. 7), it was found that the raw profiles were noisier compared to the XPORT case study. Well-structured bumps were not evidenced in the ML profiles, maybe as a result of the lower-range resolution of the MOUC radar, and the regularization procedure was found to work satisfactorily. It remains, however, difficult to assume that there is no \( \delta_{hv} \) contamination during the ML period.

Coming back to Fig. 6, one can note the mean value of \( \phi_{dp}(r_M) \) to be equal to 11.2° during the snowfall period, resulting in a specific differential phase shift on a propagation of 0.28° km\(^{-1}\) if the differential phase shift on backscatter is neglected. Such values indicate a significant heterogeneity in the horizontal and vertical dimensions of the snow/ice hydrometeors. During the rainy period between 04:10 and 06:00 UTC, there is good coherence between the specific
attenuations derived from the MRT PIA (0.078 dB km\(^{-1}\) at around 04:00 UTC–0.035 dB km\(^{-1}\) at 06:00 UTC) and those derived from the polarimetry (0.076–0.046 dB km\(^{-1}\) at the same time steps) using the \(k–K_{dp}\) relationship established for this event and by using the DSD measurements available from the IGE site (see Sect. 3.3 below).

Our main objective with the 3–4 January 2018 event is to study the \(\phi_{dp}–\text{PIA}\) relationship within the ML. Figure 6 indicates that both variables take, as expected, higher values during that period compared to during the snowfall and rainfall periods. The maximum values reached are 14.2 dB for PIA and 25.6° for \(\phi_{dp}(r_M)\). Figure 6b and c also show that the co-fluctuation of the two time series is not that good during the ML period, with a \(\phi_{dp}(r_M)\) signal having a trapezoidal shape with maximum values between 02:35 and 03:15 UTC, while the MRT PIA signal is more triangular and peaks at 03:15 UTC. We note that the two signals compare well after the peak, and that they both peak down at 03:55 UTC when measurements are made in the lowest part of the ML. These features are quite systematic for all of the 13 targets considered for the MOUC radar for this event, giving the impression that the \(\phi_{dp}–\text{PIA}\) relationship depends on the position within the ML and, as such, on the physical processes occurring during the melting. This will be further illustrated and discussed in Sect. 4.2. However, we have to mention the following three points here that may limit the validity of such inferences for the MOUC radar configuration compared to the XPORT one: (i) the MRT PIA estimates may be positively biased by radome attenuation, (ii) the polarimetry-derived PIA estimates may be affected by \(\delta_{hv}\) contamination in the ML, and (iii) nonuniform beam filling effects probably become significant for the 20–40 km range considered, leading to a smoothing of the radar signatures. There is no evidence so far of the first two points in the available dataset; this may be due to the moderate intensity of this precipitation event.

3.3 Study of the \(k–K_{dp}\) relationship in rain from in situ DSD measurements

Before presenting the analysis of the \(\phi_{dp}–\text{PIA}\) relationship in rain and in the melting layer based on the estimates for all the mountain targets and time steps available for the two sets of events, we study in this subsection the \(k–K_{dp}\) relationships that we were able to derive from the DSD measurements collected at ground level at the IGE site. For all the events, precipitation was in the form of rainfall at this altitude. As for the scattering model, we used the CANTMAT version 1.2.
software program that was developed at Colorado State University by Chenxiang Tang and Viswanathan N. Bringi. The raw Parsivel DSD measurements have a time resolution of 1 min. The volumetric concentrations were computed with a 5 min resolution and binned into 32 diameter classes with increasing sizes from 0.125 mm up to 6 mm. The CANTMAT software uses the T-matrix formulation to compute radar observables such as horizontal reflectivity, vertical reflectivity, differential reflectivity, copolar cross-correlation, specific attenuation, specific phase shift, etc., as a function of the DSD, the radar frequency, air temperature, oblateness models (e.g., Beard and Chuang, 1987; Andsager et al., 1999; Thurai and Bringi, 2005), and canting models for the rain drops and the incidence angle of the electromagnetic waves. Figure 8 displays the empirical $k$–$K_{dp}$ pairs of points obtained for the convective events (left) and the stratiform one (right) as well as the fits of least square linear models and power law nonlinear regressions.

Based on the literature review mentioning an almost linear relationship between $k$ and $K_{dp}$ at X band (Bringi and Chandrasekar, 2001; Testud et al., 2000; Schneebeli and Berne, 2012), we have first tested a linear regression with an intercept forced to be equal to 0 (red lines in Fig. 8). This simple model provides a rather good fit to the data, especially for the convective events. Due to the observed bending of the scatterplots, we have also tested a nonlinear regression to a power law model (blue curve) which significantly improves the fittings. A sensitivity analysis was performed in order to test the influence of the raindrop temperature, the raindrop oblateness model, the standard deviation of the canting angle distribution and the incidence angle. For reasonable ranges of the variation of these parameters, the DSD itself appears to be the most influential factor on the values of the regression coefficients. We note that the slopes of our zero-forced linear models are significantly higher than the values proposed in the literature (0.233 in Bringi and Chandrasekar, 2001; 0.205–0.245 in Schneebeli and Berne, 2012). The exponents of the fitted power law models are also significantly higher than 1.0. The fits in Fig. 8 correspond to the most likely parameterization of the scattering model in terms of temperature and incidence angles for the two events, i.e., 20°C and 7.5°, respectively, for the convective cases and 0°C and 0° for the stratiform case. The Beard and Chuang (1987) formulation was used as the raindrop oblateness model. The DSD-derived linear and nonlinear $k$–$K_{dp}$ relationships were used to process the regularized $\phi_{dp}(r)$ profiles which were first simply derived to obtain the $K_{dp}(r)$ profiles prior to the application of the two $k$–$K_{dp}$ relationships. The bottom graphs of Fig. 5 show examples of the resulting polarimetry-derived PIAs.

### 4 Results

#### 4.1 Study of the $\phi_{dp}$–PIA relationship in rain

Figure 9 displays the scatterplot of the $\phi_{dp}$–PIA values obtained for the nine convective events (Table 2) with the XPORT 7.5° PPI data, following the methodology described in Sect. 3.1 and 3.2. The data from the 16 mountain targets (Table 3) were considered. For a given event, targets with maximum MRT-derived PIAs less than 5 dB were discarded in order to limit the weight of small PIA estimates in the global analysis. Since we consider the two variables to be on an equal footing, we preferred to calculate the least rectangles regression (blue straight line) between the two variables rather than the least squares regression of one variable over the other one. One can notice the rather large dispersion of the scatterplot with an explained variance of 77%. We note the regression slope (0.41) to be higher than the slope of the $k$–$K_{dp}$ linear relationship (0.336), which is reported as the red straight line in Fig. 9.

To go further, Fig. 10 presents the comparison of the MRT-derived PIAs with the polarimetry-derived PIAs. The linear $k$–$K_{dp}$ relationship leads to a significant positive bias for the polarimetry-derived PIAs with a least rectangles slope of 1.24. The nonlinear $k$–$K_{dp}$ relationship does a good job of reducing this bias (least rectangles slope of 1.03). This result may be surprising given the $k$–$K_{dp}$ relationships displayed in Fig. 8. One has to realize that the range of $K_{dp}$ values is much smaller for the 5 min DSD estimations than for the $K_{dp}(r)$ profiles discretized with a 34.2 m resolution. Considering that the 1 min DSDs allowed us to confirm the validity of the linear and nonlinear $k$–$K_{dp}$ models for a wider $K_{dp}$ range (not shown here for the sake of conciseness), we are therefore confident in the relevance of the results presented in Fig. 10.

#### 4.2 Study of the $\psi_{dp}$–PIA relationship in the melting layer

Figure 11 displays the scatterplot of the $\phi_{dp}$–PIA values obtained in the ML for the 4 January 2018 stratiform event with the MOUC 0° PPI data, following the methodology described in Sect. 3.1 and 3.2. The results obtained for the 13 targets (Table 4) are considered in this analysis, with no target censoring based, for instance, on the minimum PIA observed for a given target as for the XPORT case study. One can see that the correlation between the two variables is severely degraded compared to the rain case with an explained variance of 41% and a least rectangle slope of 0.51 dB per degree. The red line recalls the $k$–$K_{dp}$ linear regression determined with the DSDs observed at ground level for this event. Clearly, the $\phi_{dp}$–PIA relationship is different in rain and in the ML and, as suggested when commenting on Fig. 6, it likely depends on the physical processes occurring during the melting.
Figure 8. DSD-derived $k$–$K_{dp}$ relationships for the nine convective events (a) and for the stratiform event of 3–4 January 2018 (b) (see text for details).

Figure 9. PIA–$\phi_{dp}$ scatterplot for the nine convective events considered in this study. The blue line corresponds to the least rectangle fit to the data, while the red line corresponds to the linear $k$–$K_{dp}$ relationship derived from the DSD data available at ground level.

To investigate this point, $\phi_{dp}(r_M)$ and PIA($r_M$) values estimated during the rising of the ML at the level of the MOUC radar are represented in Fig. 12 as a function of their position within the ML. As already noted, we define the ML width as the difference between the $Z_h$ top altitude and the $\rho_{hv}$ bottom altitude (Khanal et al., 2019). Since the ML width significantly varies during the considered period (from 630 to 1020 m; see Fig. 8), we found it necessary to scale the altitudes by the ML width. This was achieved by considering the following linear transformation of the altitudes:

$$H(t) = \left( h_M - h_{\rho_{hv}B}(t) \right) / ML_w(t),$$  \hspace{1cm} (4)

where $h_M$ is the altitude (m a.s.l.) of the MOUC radar, $h_{\rho_{hv}B}(t)$ is the altitude of the ML bottom and $ML_w(t)$ is the ML thickness at a given time $t$. The scaled altitude $H(t)$ subsequently takes the value 0 at ML bottom and the value 1 at ML top (orange and blue thick horizontal lines, respectively, in Fig. 12). Furthermore, in order to locate more precisely the position of the $Z_h$ and $\rho_{hv}$ peaks within the ML, we computed their scaled altitudes at each time step, $H_{zhP}(t)$ and $H_{\rho_{hv}P}(t)$, respectively, as follows:

$$H_{zhP}(t) = \left( h_{zhP}(t) - h_{\rho_{hv}B}(t) \right) / ML_w(t)$$  \hspace{1cm} (5)

and

$$H_{\rho_{hv}P}(t) = \left( h_{\rho_{hv}P}(t) - h_{\rho_{hv}B}(t) \right) / ML_w(t),$$  \hspace{1cm} (6)

where $h_{zhP}(t)$ and $h_{\rho_{hv}P}(t)$ are the altitudes of $Z_h$ peak and $\rho_{hv}$ peak at time $t$. The broken horizontal lines in Fig. 12 represent the 10 % and 90 % quantiles of the time series of the scaled altitudes of $Z_h$ peak (broken blue lines) and $\rho_{hv}$ peak (broken orange lines). We can observe a shift between the $Z_h$ and $\rho_{hv}$ characteristic altitudes, consistent with the ML climatology established by Khanal et al. (2019) who reported a shift of about 100 m in average between the two peaks. We note in Fig. 8 that this shift is visible during the snowfall period and at the beginning of the ML rising but that it is less pronounced after 03:00 UTC and during the rainfall period. In order to better evidence their vertical trends, the MRT PIA($r_M$) and $\phi_{dp}(r_M)$ values are presented in Fig. 12 as a function of the scaled altitudes in the form of boxplots with a scaled altitude class of size 0.1. The number of counts in each
class is indicated on the right of the graphs; it is a multiple of the number of MRT targets (13 here) depending on the time occurrence of estimates in a given altitude class. The vertical sampling is not very rich, with missing classes within the ML. However, there is a clear signature for the two variables in the ML. The trends already evoked when commenting on Fig. 8 are confirmed as follows: (i) the MRT PIAs peak when measurements are made at the level of the $Z_h$ and $\rho_{hv}$ peaks – more precisely, the PIA peak is observed for the altitude class containing the $\rho_{hv}$ peaks (scaled altitude class centered at 0.3); (ii) the region with maximum values is somewhat thicker for $\phi_{dp}$, encompassing a significant part of the upper ML, between the 0.3 and 0.8 scaled altitude classes; (iii) $\phi_{dp}$ tends towards almost similar values on average in rain (ML bottom) and snow (ML top); and (iv) the PIA tends towards its value in rain below the ML and towards 0 above the ML. One would have expected a more pronounced return towards 0 of the PIAs on top of the ML. This lower-than-expected decrease could sign a radome attenuation; however, the rainfall intensity is low for the considered event, and the radome is equipped with a heating system so that accumulated snow is unlikely. It may also result from a smoothing effect related to nonuniform beam filling. With its 3 dB beamwidth of 1.28°, the angular resolution of the measurements of the MOUC radar is 447 and 1005 m at distances of 20 and 45 km, respectively, which correspond to the minimum and maximum ranges of the considered mountain targets.

Finally, Fig. 13 displays the evolution of the ratio of the mean of the MRT PIA ($r_M$) values over the mean of $\phi_{dp}$ ($r_M$) values as a function of the scaled altitudes. The value of the ratio below the ML (0.33) is in rather good agreement with the slope of the linear model established between the specific attenuation $k$ and the specific differential phase shift $K_{dp}$ using the DSD measurements in rain available for this event (0.29; see Fig. 9). Near the $\rho_{hv}$ peak, the ratio value is equal to 0.42. For the three classes of scaled altitude 0.7, 0.8 and 0.9, the ratio is between 0.32 and 0.38, with an apparent secondary maximum for the altitude class 0.8. Data with increased vertical resolution would be necessary to confirm, or not, this observation, which is also visible on the PIA profile and on several $\phi_{dp}$ and PIA time series like the ones displayed in Fig. 8. Above the ML, the ratio progressively tends toward 0 at about 300 to 400 m.

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Boxplots of the PIA and $\varphi_{dp}$ values within the ML as a function of the scaled altitude (a and b, respectively) for the stratiform event of 4 January 2018. The horizontal blue and orange continuous lines represent the ML top and bottom, respectively; the broken horizontal blue and orange lines give the 10% and 90% quantiles of the scaled altitudes of the $Z_h$ and $\rho_{hv}$ peak distributions, respectively.

Figure 13. Evolution of the ratio of the mean PIAs over the mean $\varphi_{dp}$ values within the ML as a function of the scaled altitudes for the stratiform event of 3–4 January 2018. The horizontal blue and orange lines represent the ML top and bottom, respectively; the broken horizontal blue and orange lines give the 10% and 90% quantiles of the scaled altitudes of the $Z_h$ and $\rho_{hv}$ peak distributions, respectively.

5 Summary and conclusions

In this work we developed a methodology for studying the relationship between total differential phase shift ($\varphi_{dp}$) and path-integrated attenuation (PIA) at X band. Knowledge of this relationship is critical for the implementation of attenuation corrections based on polarimetry. We used the mountain reference technique (MRT) for direct PIA estimations associated with the decrease in strong mountain returns during precipitation events. The MRT sensitivity depends on the time variability of the dry weather mountain returns. The MRT PIAs may be positively biased by on-site attenuation related in particular to radome attenuation and negatively biased by the effect of precipitation falling over the reference targets. The polarimetry PIA estimation is based on the regularization of the raw $\varphi_{dp}$ profiles, and their derivation in terms of specific differential phase shift ($K_{dp}$) profiles, followed by the application of a power law relationship between the specific attenuation and the specific differential phase shift. Such $k$–$K_{dp}$ relationships were evaluated for rain with a scattering model by using DSD measurements and an oblateness model for raindrops. The noise of the raw $\varphi_{dp}$ profiles, the possible contamination of the signal by differential shift on backscatter and the adequacy of the $k$–$K_{dp}$ relationship is the main factor that determines the quality of the polarimetry-derived PIAs. Nonuniform beam filling (NUBF) effects may also play a role. A point to emphasize is that both PIA estimators are not sensitive to an eventual radar miscalibration.

We presented first a rain case study based on nine convective events observed with the XPORT radar located in the Grenoble valley. A total of 16 mountain targets were considered with the dry weather mean apparent reflectivity greater than 45 dBZ. The stability of the apparent reflectivity of the mountain targets was shown to be very good, which is an indication of good radar calibration stability during the considered period. The time variability of the reference returns during dry weather preceding or succeeding the rain events was also found to be very small with standard deviations in the range of [0.2–0.9 dBZ], enabling a MRT PIA sensitivity better than 1 dB. Since the XPORT radar is radomeless, on-site attenuation effects are most likely negligible. The im-
pact of rain falling over the mountain targets may also be very limited due to the high reflectivity threshold considered (45 dBZ). The development of the regularization procedure of the raw $\psi_{dp}$ profiles required a significant effort, and we are confident in its ability to deal with the measurement noise, especially for heavy precipitation. We carefully examined many raw and regularized profiles, looking for possible evidence of $\delta_{hv}$ contamination during the considered convective events. We found some profiles with rather well organized bumps that could signal such contaminations. The regularization procedure was adapted in order to filter such effects, with a satisfactory performance when they occur at some distance (some kilometers) from the mountain target. In addition, we remind the reader that the observed $\psi_{dp}(r_M)$ values extend up to $80^\circ$, while the theoretical $\delta_{hv}$ range is 0–4 dB. The $\delta_{hv}$ effect may therefore impact the results obtained only at the margin in the considered case study. NUBF effects may constitute an additional source of error which, although the rain events were convective, should remain limited due to the short ranges considered. In the end, the scatterplot of the MRT PIA as a function of the $\phi_{dp}(r_M)$ values for all the nine convective events presents a good coherence overall with, however, a significant dispersion (explained variance of 77 %). It is interesting to note that the nonlinear $k$–$K_{dp}$ relationship derived from independent DSD measurements taken during the events of interest at ground level allows for a satisfactory transformation of the XPORT $\phi_{dp}(r_M)$ values into almost unbiased (although dispersed) PIA estimates. Both estimation methods are prone to specific errors and, even if the MRT PIA estimator is more directly related to power attenuation, it is a priori difficult to say which estimator is the best. An assessment exercise of attenuation correction algorithms, making use of both PIA estimators, with respect to an independent data source (e.g., rain gauge measurements), is desirable to distinguish the two PIA estimators. From this perspective, a specific experiment is being designed within the RadAlp project and it will be implemented in the near future.

The melting layer (ML) case study of 3–4 January 2018 was made possible by the unique configuration of the observation system available. The study of the $k$–$K_{dp}$ relationship within the ML is desirable to better quantify the attenuation effects in the ML with polarimetry; and one has to recognize that such a relationship can still be very difficult to characterize theoretically given that scattering models and particle size distributions need to be collected in the ML. The XPORT radar located at the bottom of the valley allowed for a detailed temporal tracking of the ML from below using quasi-vertical profiles derived from 25° PPIs. The MOUC radar provided horizontal scans at an altitude of 1917 m a.s.l. in the direction of several mountain targets during the rising of the ML in about 2 h. From this dataset, it was possible to derive the evolution of PIA ($r_M$) and $\phi_{dp}(r_M)$ values as a function of the altitude within the ML. The evolution with the altitude of the ratio of the mean value of PIA ($r_M$) over the mean value of $\phi_{dp}(r_M)$, as a proxy for the slope of a linear $k$–$K_{dp}$ relationship within the ML, was also considered. Since the ML width varied during the ML rising, we found it necessary to scale the altitudes with respect to the ML width. The three variables considered present a clear signature as a function of the scaled altitude. In particular, the PIA/$\phi_{dp}$ ratio peaks at the level of the $\rho_{hv}$ peak (somewhat lower than the $Z_h$ peak), with a value of 0.42 dB per degree, while its value in rain just below the ML is 0.33 dB per degree. The latter value is consistent with the slope of the linear $k$–$K_{dp}$ relationship (0.29) established from concomitant DSD measurements at ground level. The PIA/$\phi_{dp}$ ratio remains quite strong in the upper part of the ML, between 0.32 and 0.38 dB per degree, before tending towards 0 above the ML. One would have expected a more pronounced return towards 0 of the PIA on top of the ML. This lower-than-expected decrease could signal on-site attenuation occurring at the beginning of the ML rise due to the melting of the snow eventually accumulated over the radome; this effect is probably low for the considered event since the snowfall intensity was small and since the radome is heated. It may also result from a smoothing effect related to nonuniform beam fanning (angular resolution of 447 and 1005 m for the range of mountain target distances). The $\delta_{hv}$ effect is likely to be strong in the ML (up to 4°), and its relative importance may be quite high in our case study since the PIA range is significantly lower compared to the rain case study, with maximum PIA of about 15 dB (note also that the sensitivity of the MRT is less than for the XPORT case study since the dry weather variability of the mountain returns is higher with standard deviations in the range [0.62–1.44]). However, we did not find evidence of $\delta_{hv}$ signatures in the raw $\psi_{dp}(r)$ profiles, and we are confident in the ability of the regularization procedure to filter them in a rather satisfactory way if they eventually occur. Although the experimental configuration for the study of attenuation in the ML presents some limitations (e.g., possible radome attenuation and NUBF effects), the preliminary results presented here will be deepened by processing a dataset of about 30 stratiform events with the presence of the ML at the level of the MOUC radar.

Data availability. Data can be made available by the first author upon request.

Author contributions. GD developed the concept of the article, realized the codes and performed the data processing; GD also wrote the article and handled the review process. AKK was involved in the development of the codes (ML identification, $\phi$dp processing, etc.); NY shared his experience of the RHytMME radars and provided the MOUC radar data and the preliminary version of the $\phi$dp regularization procedure. FC is a research engineer who has built the XPORT radar and is now responsible for the current updating program and the data acquisition. BB and NG contributed to the article through scientific discussions and amendments to the paper.

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