Heterometallic 3d–4f single-molecule magnets

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The promising potential applications, such as information processing and storage or molecular spintronics, of single-molecule magnets (SMMs) have spurred on the research of new, better SMMs. In this context, lanthanide ions have been seen as ideal candidates for new heterometallic transition metal–lanthanide SMMs. This perspective reviews 3d–4f SMMs up to 2014 and highlights the most significant advances and challenges of the field.

Introduction

In the early 1990s the discovery that transition metal complexes could retain magnetization caused an explosion in the field of coordination complexes. The first molecule that was shown to act as a magnet at the molecular level was a dodecanuclear complex of \( \text{Mn(II)}/\text{Mn(IV)} \) with acetato, oxo and water ligands, \([\text{Mn}_{12}\text{O}_{12}(\text{MeCOO})_{16}(\text{H}_2\text{O})_4]\), \( \text{Mn}_{12}\text{Ac} \). Molecules with this property were then called single-molecule magnets (SMMs). Soon other transition metal complexes were also found to be SMMs, including a large family of \( \text{Mn}_{12} \) complexes with structures related to that of \( \text{Mn}_{12}\text{Ac} \). \( \text{Mn}_{12}\text{Ac} \) was the first SMM discovered and is still one of the most studied. SMMs are by themselves already very interesting molecules with very special magnetic properties, but their discovery also brought about the possibility of their use in technological applications, substituting conventional ferromagnetic materials for SMMs. The promise of the use of SMMs in the processing and storage of information is not only that of the use of a new material for the same task, but that of opening up the possibility of having ultra-high density information storage devices, or ultra-fast information processing devices based on SMMs. Additionally, as new SMMs are discovered, new applications are proposed, for example their use in molecular spintronics. There are two main problems that must be solved before all these proposed technological applications of SMMs can be implemented. Firstly, the working temperatures must be improved, so far all SMMs discovered to date only function at liquid helium temperatures. This in itself does not make them useless, but it greatly holds them back from any applications in information storage and processing, which is nowadays efficiently done with bulk ferromagnetic materials, or even spintronics. Secondly, in order to fabricate devices using SMMs, new technologies must be developed; so far there are only a few examples of SMMs for which surface deposition and addressing of single molecules has been explored. One can envision temperature not being an issue if the physical properties of the new SMMs should out-perform the classical magnets used today. Still, the issue of depositing and addressing a single molecule on a surface remains a great challenge for scientists working in the field.

The promise of the ultimate miniaturization of information storage and processing devices using SMMs has driven many researchers’ efforts in obtaining new improved examples of SMMs. In particular, the main goal has been to obtain SMMs with higher working temperatures, but there are still many different ways in which researchers report the success of new SMMs; a useful parameter would be the blocking temperature. The blocking temperature, \( T_b \), the maximum temperature at which the SMM is functional, should be the temperature at which magnetization hysteresis vs. field is observed. The use of the blocking temperature is greatly hampered due to the fact that its value greatly depends on the sweep rate of the magnetic field during the measurement and on the experiment used to measure it, so when comparing blocking temperature values one should be extremely careful. When characterizing SMMs, the effective barrier for reversal of the magnetization, \( U_{\text{eff}} \), is most often reported in the literature. This is also called the anisotropy barrier and it is the energy needed to transform the SMM into a simple paramagnet. \( U_{\text{eff}} \) is the most popular parameter used to characterise SMMs, mainly due the phenomenon of quantum tunnelling of the magnetization, of particular importance for 3d–4f SMMs. For a complex to be a good SMM with a high blocking temperature, \( U_{\text{eff}} \) must be large. Several researchers have proposed ways in which to normalise the parameters that should be reported for each new SMM. Long and co-workers proposed that the temperature at which 1/2 width hysteresis is observed should be reported, while Sessoli and co-workers proposed to define the
Blocking temperature as the temperature at which the magnetisation relaxes in 100 s. However, neither of these two definitions is widely used by scientists in the field and \( U_{\text{eff}} \) is still the parameter that is usually reported. In part, this is due to the fact that most reported SMMs do not display hysteresis of the magnetization above 1.8 K, the lower limit in temperature for commercial SQUID magnetometers, thus access to lower temperatures is required to perform magnetisation vs. field hysteresis measurements to obtain \( T_b \). The \( Mn_{12} \) family of SMMs, with an \( S = 10 \) ground state, were the SMMs with the highest blocking temperatures \( (T_b = 3.5 \text{ K}) \) and \( U_{\text{eff}} \) values up to 74 K until 2007, when Brehcin and co-workers reported a \( Mn_6 \) complex with a record \( U_{\text{eff}} \) of 86.4 K and \( S = 12.5 \). Thus, many SMMs have been reported since \( Mn_{12}Ac \), but reported working temperatures still remain in the liquid helium range. If anisotropy barriers are compared the case is similar, and the reported values are not that much greater than that of \( Mn_{12}Ac \).

If controlling spin is a hard task for experimental chemists, controlling the anisotropy of a high nuclearity complex is even more complicated. Initially, complexes of the more anisotropic transition metals, known to display strong axial anisotropy, such as \( Mn^{\text{III}} \), \( Co^{(ii)} \) and \( Ni^{(ii)} \), were investigated, and still are, as a means to obtain better SMMs. However, in the early 2000s the search for 3d–4f SMMs became an important trend, with the goal of improving the anisotropy of the obtained species and thus obtaining better SMM properties. The lanthanide ions are well known for having strong spin–orbit coupling, their magnetic properties ruled by the quantum number, \( J \), which has the maximum value of \( |L + S| \) for lanthanides with more than half-filled f shells \( (Tb, Dy, Ho, Er, Tm, Yb, Lu) \) and the minimum value of \( |L - S| \) for lanthanides with less than half-filled f shells \( (Ce, Pr, Nd, Pm, Sm, Eu) \). The number of unpaired electrons has little impact on the magnetic moment of the lanthanide ion and those with large \( m_J \) values of the ground state are the ones with stronger magnetic moments. The ground state bistability characteristic of an SMM arises in lanthanide ions from the \( m_J \) sublevels of the \( ^{2S+1}L_J \) term. The most common lanthanide ions used to obtain SMMs are terbium(III) and dysprosium(III), but also erbium(III), samarium(III), ytterbium(III), gadolinium(III) and holmium(III) have been used. As Ishikawa showed for the Tb and Dy sandwich phthalocyanin complexes reported in 2003,8 in lanthanide SMMs the energy barrier is defined by the spin and angular momentum of a single lanthanide placed in a ligand field giving the largest \( |J_z| \), the lowest energy, and a large energy gap to the next sublevels. Long and Rinehart proposed simple rules in order to exploit the lanthanides’ single-ion anisotropy for designing 4f SMMs.9 According to their theory, to maximize the anisotropy of oblate ions \( (Ce^{(iii)}, Pr^{(iii)}, Nd^{(iii)}, Tb^{(iii)}, Dy^{(iii)} \) and \( Ho^{(iii)} \)) the crystal field should be such that the ligand electrons are concentrated above and below the \( xy \) plane. On the other hand, for prolate ions \( (Pm^{(iii)}, Sm^{(iii)}, Er^{(iii)}, Tm^{(iii)} \) and \( Yb^{(iii)} \)) an equatorial coordination geometry is preferred. Many of the reported 4f SMMs follow this prediction, and it is particularly useful for mononuclear lanthanide SMMs. This simple qualitative way of predicting SMM behaviour could also be used to ascertain whether a lanthanide ion in a 3d–4f complex will contribute strongly to the complex anisotropy, and thus, to the SMM properties of the 3d–4f species. However, this must only be considered in a very qualitative manner. Given the difficulties in factoring out all of the contributions to a polynuclear complex’s magnetic anisotropy, the relationship between the ligand arrangement around the

3d–4f SMMs: exploiting the magnetic properties of the 4f ions

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lanthanide ion in a 3d–4f polynuclear complex and the complex’s axial anisotropy will not be as straightforward as with mononuclear lanthanide SMMs. Ideally, the most anisotropic 3d metals should be combined with the right lanthanide to obtain new SMMs: the anisotropy of the 3d–4f complex will be a combination of the single-ion anisotropies of all the paramagnetic metal centres involved. A huge advantage of the lanthanide ions is the possibility of preparing families of complexes in which the properties of the 3d metal core can be checked, for example using the diamagnetic Y(III) or the lanthanides La(III) or Lu(III), with similar ionic radii, or the isotropic Gd(III). In this way, one can check the contribution to the SMM property from the 3d metals in the molecule; for example, the family of complexes 5 in Table 1, where the 5 (La) without 4f electrons or anisotropy provided by the lanthanide ion is an SMM, indicating that the [Mn₆] core of the molecule is the main contribution to the anisotropy. Using a diamagnetic 3d metal analogue might be feasible, as many researchers have done in the past for diatomic oxo bridges are the surest way to enforce the strongest possible coupling. Usually, 3d–4f exchange constants have values below 5 K. A great tool to elucidate 3d–4f magnetic coupling is the software

Table 1 3d–4f SMMs. The effective barriers measured with an applied dc field have the field value in Oe in parentheses

| Formula | U_eff (K) | T_0 (K) | Ref. |
|---------|-----------|---------|-----|
| [CuLn(hfac)_2](Ln = Tb, Dy) | 2.1 | 1.2 | 19 |
| Dy₂Mn₆(H₂Hsh)[HshH][MeOH] | 20 |
| Mn₁₁Dy₂O₆(OH)₂(OMe)₃(O₂CPh)_4(NO₃)₆(H₂O)_2 | 9.3 | 21 |
| Mn₈Dy₂O₄Cl₂(N₆H₄)₄(H₂O)₂ | 0.5 | 23 |
| Pr₄N⁺[Mn₄Lu₄(OMe)₃(SALtO)₂(SALtOH)] | 6; 1.3 | 10 |
| M₇[Lu]₂[PhCOO]Dy₁₂[hfac] | M = Zn, Co | 47.9; 8.8 | 26 |
| Co₂Dy[OECm][teaH]₄(Piv)₆ | 51; 127 | 29 |
| Dy₂Co₂[OECm][L₃OOCPh] | NO₃ | 87, 104 | 31 |
| Dy₂Co₂[OECm][L₃OOCPh] | NO₃ | 79, 115 | 31 |
| Dy₂Co₂[OECm][teaH][OECm][MeOH] | NO₃ | 88 | 32 |
| Cr₂Dy[OECm][R]d[Dy]₄[OECm] | NO₃ | 34, 37, 41 | 1.8, 2.2, 2.2 | 33 |
| Cr₂Dy[OECm][O₂CPh][mdea][NO₃] | 77 | 2.2 | 33 |
| Dy₆Mn₉[OECm][CymCOO]₈(Mpea)₈(mp)₂(MeOH)₂(H₂O)₂ | 19.6; 47 |
| Ni₂Dyₓ[OECm][S] | S = MeOH, DMF | 21.3; 18.5 | 4.2; 3.2 | 35 |
| Co₂Dy[OECm][NO₃] | 82 | 3 | 36 |
| MnₓLuₓ(OECm)[mdeaH][Piv][NO₃] | NO₃ | 38.6 | 1.9 | 37 |
| Mn₁₃Dy₂O₆(OH)₂[BuCOO]₄(HCOO)[NO₃] | NO₃ | 74 | 3 | 38 |
| Dy₂Co₂[CO₃]₄(OCN)₄(MNO₃)₂(Ph₃CO)₄(MeOH) | 4.3; 25 | 39 |
| H₂O[Co₂Dy₂[OECm][CymCOO]₈(Mpea)[MeOH] | NO₃ | 4.6 | 0.6 | 40 |
| MnₓDyₓ[OECm][S] | S = MeOH, DMF | 103; 32.8 | 3.1; 8.7 | 45 |
| MnₓDyₓ[OECm][S] | S = MeOH; Dy = Tb, S = 3 MeOH 1 EtOH | 24; 46 | 46 |
| MnₓDyₓ[OECm][S] | S = MeOH; Dy = Tb, S = 3 MeOH 1 EtOH | 36.6 | 47 |
| TbxMn₂[OECm][mdeaH][Ph₃CO][MeOH] | NO₃ | 19.6 | 47 |
| TbxMn₂[OECm][mdeaH][Ph₃CO][MeOH] | NO₃ | 33 | 37 |
| MnₓDyₓ[OECm][S] | S = MeOH; Dy = Tb, S = 3 MeOH 1 EtOH | 13.83 | 48 |
| MnₓDyₓ[OECm][S] | S = MeOH; Dy = Tb, S = 3 MeOH 1 EtOH | 25 | 49 |
| (CuL)₂[MeOH][NO₃] | 20.3 | 50 | 50 |
| (CuL)₂[MeOH][NO₃] | 18 | 50 | 50 |
| (CuL)₂[MeOH][NO₃] | [1] | 19.3 (1000) | 1.6 | 52 |
| (CuL)₂[MeOH][NO₃] | [2] | 19 | 53 |
| (CuL)₂[MeOH][NO₃] | 13 | 54 |
| (CuL)₂[MeOH][NO₃] | 15.6 | 55 |
| (CuL)₂[MeOH][NO₃] | 13 | 0.7 | 56 |
| (CuL)₂[MeOH][NO₃] | 30 | 57 |
| (CuL)₂[MeOH][NO₃] | 15 | 58 |
| (CuL)₂[MeOH][NO₃] | 16 | 0.5 | 59 |
| (CuL)₂[MeOH][NO₃] | 12 (2000); 12 | 1 | 61 |
| (CuL)₂[MeOH][NO₃] | 19 (1000) | 70 |
| (CuL)₂[MeOH][NO₃] | 27 (1000) | 70 |
| (CuL)₂[MeOH][NO₃] | 19, 92 | 73 |
behaviour. A typical example is shown in Fig. 1. These experimental frequency-dependent is the most reliable signature of SMM

...the ac oscillating field will freeze and lag behind the applied ac field resulting in a complex. However, that would not be the case for Tb(III) 3d

As for any SMM, 3d–4f SMMs are usually characterized in their pure crystalline form using commercial SQUID magnetometers. The usual measurements of susceptibility against temperature and magnetization against field are also performed for 3d–4f SMMs. With these data, usually reported as the χ′T product and the reduced magnetisation, one can evaluate the magnetic coupling between the metal centres in the complex as well as the spin ground state. For 3d–4f complexes, this is by no means straightforward and in many it will not be possible to quantify the magnetic exchange. For SMMs, alternate current (ac) magnetic susceptibility is also measured. Usually a small magnetic field of 1–5 Oe that oscillates at frequencies between 1 and 1500 Hz is used to measure magnetic susceptibility over a range of temperatures. For an SMM each individual molecule has an energy barrier to be overcome in order to reverse the magnetic moment. The SMM molecules will freeze and lag behind the applied ac field resulting in a susceptibility signal with two components: one in-phase with the ac oscillating field χ′, and one out-of-phase with the oscillating field χ″. The appearance of out-of-phase maxima that are frequency-dependent is the most reliable signature of SMM behaviour. A typical example is shown in Fig. 1. These experiments are usually performed by scanning the temperature and frequency domains, resulting in susceptibility vs. T and susceptibility vs. frequency plots. From the slope of the plot of lnτ vs. 1/T one can obtain the anisotropy barrier where the graph is linear using an Arrhenius type of equation, τ ≈ τ0exp(Ueff/kT). For lanthanide SMMs, a temperature-independent region is usually observed that is characteristic of fast relaxation of the magnetisation via quantum tunnelling (QTM). When the thermal mechanism coexists with the QTM, a curvature is seen in the Arrhenius plot of lnτ vs. 1/T. Several thermal relaxation processes can also coexist, and this can be assessed by examination of so-called Argand or Cole–Cole plots. In these plots, χ″ is plotted against χ′ at a constant temperature, resulting in a semi-circular representation. When a distribution of energy barriers exists the semicircle is distorted, a Debye model applies and the parameter a gauges the distribution. Usually for SMMs this parameter has values smaller than 0.2.

Researchers still have a lot to learn from new lanthanide SMMs,25 including the occurrence of multiple relaxation processes for pure SMMs,71 and the toroidal magnetic moments in some 4f SMMs,72 among other things. The phenomena of multiple relaxation processes and QTM also occur in transition metal SMMs.

Magnetisation vs. field hysteresis loops can be measured in a commercial SQUID magnetometer down to 1.8 K or using a micro-SQUID or micro-Hall probe with suitable sweep rates for the field, going down to the mK range. Due to the fact that most 3d–4f SMMs display hysteresis of the magnetisation vs. field at very low temperatures, when possible micro-SQUID data are reported. Blocking temperatures and anisotropy barriers can be obtained from magnetisation decay measurements, also performed using a micro-SQUID. X-ray magnetic circular dichroism (XMCD) is a technique that has been used on a few occasions to analyse the magnetic behaviour of 3d–4f complexes. XMCD is element-sensitive and can be used to probe the magnetism of each metal type in a heterometallic complex. Pedersen and co-workers published a great example in 2012, where they reported element-specific curves for a Cr–Dy complex.13 XMCD could be exploited for evaluating the 3d–4f coupling in heterometallic complexes. Furthermore, XMCD can also be used to probe the magnetic properties of an SMM on a surface, a key and challenging step for physicists and chemists, as we reported for a Dy MOF of SMMs,14 and as Sessoli and co-workers pioneered with their work on Mn12,15 and Fe416,17 SMMs on surfaces.18

A wide perspective on 3d–4f SMMs

The first 3d–4f SMMs were reported in 2004: [CuLLn(hfac)2]2,

(1, Ln = Tb, Dy; H3L = 1-(2-hydroxybenzamido)-2-(2-hydroxy-3-methoxy-benzylideneamino)-ethane)19 squares with the 3d and 4f metals in an alternated array reported by Matsumoto et al., and the [Dy2Mn4(H3shi)2(Hshi)2(Hshi)10(MeOH)10(H2O)2]2,

(2, H3shi = salicylhydroxamic acid) complex, reported by Pecoraro.20

![Fig. 1 Typical ac magnetic susceptibility data and Arrhenius plot for a 3d–4f SMM. Reproduced from ref. 46 with permission from the Royal Society of Chemistry.](image)
These complexes presented no hysteresis down to 1.8 K, the temperature limit of commercial SQUID magnetometers, but ac magnetic susceptibility studies showed out-of-phase signals, as expected for SMM behaviour. Only for 1(Tb), [CuLTb(hfac)_2], were the maxima of the out-of-phase signals observed above 1.8 K. In the same year, Christou and co-workers reported the first observed magnetization vs. field hysteresis loops for a heterometallic 3d–4f SMM,[21] the [Mn_{11}Dy_{4}O_{8}(OH)_{6}(OMe)_{6}O_{2}CPh]_{6}([NO_{3})_{5}]_{2}[H_{2}O]_{3}] complex, 3. The out-of-phase ac magnetic susceptibility data showed no maxima down to 1.8 K, but magnetization vs. field hysteresis loops were observed using a micro-SQUID device below 1 K. The hysteresis loops were smooth, without the presence of the typical QTM steps signature of transition metal SMMs. This lack of steps in the magnetization vs. field hysteresis was attributed to a distribution of molecular environments or intermolecular interactions, both factors known to smooth or smear hysteresis loops in some transition metal SMMs.[22]

From these three first reports one can already make a clear conclusion: the ligands used are not specific for lanthanide–3d complexes, but are the usual varied ligands used in transition metal chemistry and lanthanide chemistry. Also from only these three reports one can see how in complex 2 the dysprosium ions are close to each other and can be coupled, while in 1 and 3 the lanthanide ions are separated by the transition metals and could only couple to the transition metals.

The first results obtained seemed encouraging and efforts were doubled. In 2004 three peer reviewed research papers reported 3d–4f SMMs. From 2008 the number increased to between 10 and 16 papers each year reporting new 3d–4f SMMs. In 2014 the exponential growth in the field was reflected in more than 50 peer reviewed research papers reporting new 3d–4f SMMs. To the date of this report, 161 3d–4f SMMs have been reported. Nearly 2/3 of the reported 3d–4f SMMs contain Dy(m) as the lanthanide ion, and in most cases the coordination environment can be described as ligands above and below the plane, thus providing the ideal setting for SMM properties for an oblate ion like Dy(m). One case of a neodymium 3d–4f complex, the [Mn_{4}Nd_{3}] SMM, is reported, but no hysteresis was observed.[24] The 3d metal is usually Co(m) or Mn(m), both highly anisotropic, but also Fe(m), Ni(m), Cu(m) or even the diamagnetic Co(m) and Zn(m) have been used. Worth mentioning here are the reported 3d–4f SMMs that contain Gd(m) and La(m). In these two cases the lanthanide is either isotropic (Gd[3m]) or diamagnetic (La[3m]). All the 3d–4f SMMs reported with these two lanthanide ions are clearly cases where the anisotropy and the SMM property are both provided for by the 3d metal part of the complex.

By design: metal substitution

As with any polynuclear coordination complex, the synthesis of 3d–4f SMMs most often follows a procedure of serendipitous self-assembly, where researchers try to provide the best reaction conditions to obtain complexes that might be new examples of SMMs. This is why there is such a rich structural diversity of 3d–4f SMMs, as is the case for transition metal SMMs. Of course the counterpart is the lack of control in the structure and properties of the prepared complexes. In the last few years the targeted substitution of a 3d metal by a lanthanide ion in a known transition metal polynuclear complex has been successfully done. This method has led to the isolation of 3d–4f complexes, where the position of the lanthanide ion could be predicted at the synthesis step. The first example was reported by Powell and co-workers in 2009,[23] when they succeeded in replacing the central Mn(n) atom of a ferromagnetically coupled [Mn_{10}] complex, [Mn_{10}O_{6}(N_{3})_{6}(HL)_{12}(MeCN)_{5}]_{2}Cl_{2}, with no anisotropy, for a Dy(m), to obtain [Mn_{18}DyO_{4}(Cl)_{6}][N_{3}]_{1.5}(HL)_{13}(MeOH)_{6}Cl_{4}, (4). By this replacement, retaining the core topology of the cluster, the anisotropy of the complex was enhanced and the SMM property observed. Thus, the introduction of the anisotropic Dy(m) ion results in the onset of the SMM property, which was absent in [Mn_{10}]. Powell’s complex, 4 (Mn_{18}Dy), provided a sandwich type of ligand environment to the Dy(m) ion and thus complex 4 displayed SMM properties. [Mn_{10}] had a record S = 83/2 spin but lacked any appreciable anisotropy. The introduction of the Dy(m) ion in a sandwich-like crystal field provided the necessary anisotropy to observe SMM behaviour.

More recently we reported a [Mn_{11}] species,[24] with a complex structure that was formed by three [Mn_{9}], units centred around a Mn(m) ion in a very large cavity for a transition metal. Conditions were perfect for the controlled preparation of a Mn–Ln complex. We tweaked the reaction conditions and we successfully isolated the desired complexes [PrNH_{2}]_{3}[Mn_{9}LnO_{6}(OMe)_{6}(SALO)_{6}(SALOH)_{3}] with several lanthanide ions: La, Gd, Tb and Dy (5).[25] The introduction of the lanthanide ion resulted in slightly enhanced SMM properties in 5(La) and modified SMM properties for 5(Dy) and 5(Tb). In 5(Gd) the SMM properties were completely quenched, probably due to better magnetic coupling through the Gd(m) ion. In a qualitative, simple manner, this can be explained by looking at the coordination environment of the Ln(m) ion, which in the complexes 5(Ln) was highly symmetrical, resembling a spherical arrangement of the ligands around the lanthanide, which according to Long’s qualitative considerations should not provide a good crystal field for a bistable ground state for Dy(m) or Tb(m). In order to substitute a 3d metal for a lanthanide ion in a known complex there must be a metal site that is appropriate for the lanthanide. This is not straightforward and there are not many examples in the literature where controlled substitution is reported.

Dy(m) 3d–4f SMMs

Without a doubt dysprosium is the most used lanthanide in order to obtain 3d–4f SMMs. The first reports in 2004 were of 3d–Dy SMMs. 3d–Dy SMMs have been reported for cobalt, chromium, copper, iron, manganese, nickel and zinc.

In several of these examples, the 3d metals in these species are Co(m) and Zn(m), thus diamagnetic, and it can be argued whether these should be considered 3d–4f SMMs or simply lanthanide SMMs with metalloligands. We have decided to consider these species here since usually when lanthanide...
SMMs are reviewed these species are not accounted for.\(^2,2^5\) Chaudhury and co-workers propose to use diamagnetic Zn(i) ions in [M\(_2\)(L)\(_2\)(PhCOO)\(_2\)Dy\(_2\)(Hfac)]\(_4\) (6(Zn) and 6(Co) in Table 1) to enhance the energy barrier of Dy(III) SMMs, as they show with DFT and \textit{ab initio} calculations. This is also studied by Shanmugam and co-workers.\(^2,2^6\) The record anisotropy barrier for 3d–4f SMMs belongs to a family of structurally related [Co(III)\(_2\)Dy(III)] complexes with a double defective cubane structure prepared using triposadal ligands, the metal–oxo core is shown in Fig. 2. This core topology is well known in transition metal chemistry, and many Mn\(_{ni}\)Mn\(_{ni}\) complexes have been reported with this core that are SMMs, a review of their properties and structures can be found in Aromi and Brechin’s review.\(^2,2^8\) The highest reported anisotropy barrier for a 3d–4f SMM belongs to a family of complexes, with a Type A defective cubane core, as shown in Fig. 2. [Co\(_2\)Dy\(_2\)(OMe)\(_2\)(teaH)\(_2\)(Piv)\(_6\)] (7, teaH = triethanolamine, Piv = pivalate)\(^29\) shows two overlapped peaks in its out-of-phase ac magnetic susceptibility, with \(U_{eff} = 51\) and 127 K, as shown in Fig. 3. In 2013 and in 2014, Murray and co-workers reported two families of [Co(III)\(_2\)Dy(III)] (8 and 9 in Table 1) with Type A metal arrangements, as shown in Fig. 2.\(^2,3^0\) This work added to their previous report of a [Ln\(_2\)Co\(_2\)(OMe)\(_2\)(teaH)\(_2\)(OOCPh)\(_4\) (MeOH)]\(_2\)(NO\(_3\))\(_2\)] family of complexes in 2012\(^3^2\) with the same defective cubane core. The dysprosium–cobalt analogue, [Dy\(_2\)Co\(_2\)(OMe)\(_2\)(teaH)\(_2\)(OOCPh)\(_4\)(MeOH)]\(_2\)(NO\(_3\))\(_2\)] (complex 10 in Table 1), had a record energy barrier at the time of 88 K, and its ac data could be fit to a distribution of energy barriers with \(\alpha = 0.25\), attributed by the authors to the fact that there are two independent Dy(iii) sites.

The authors were able to evaluate the magnetic exchange between Dy(III) ions, which was weak and antiferromagnetic. QTM was suppressed in a bulk sample of complex 10 due to the weak dipolar antiferromagnetic coupling, but dilution experiments in a yttrium(III) analogue matrix showed that fast tunnelling dominated the magnetic relaxation. The same authors showed how the structurally related [Cr\(_2\)Dy\(_2\)(OMe)\(_2\)–(Rdea)\(_2\)(acac)\(_2\)NO\(_3\))\(_2\)] (Rdea = R-diaethanolamine; 11(Me) R = methyl, 11(Et) R = ethyl or 11(Bu) R = tert-butyl) and [Cr\(_2\)Dy\(_2\)(OMe)\(_2\)(mdea)\(_2\)(O\(_2\)CPh)\(_2\)](NO\(_3\))\(_2\)] (12), with a Type A defective cubane core as shown in Fig. 2, had large anisotropy barriers and long relaxation times compared to the Co(iii) analogues due to the significant magnetic interaction between the 3d metal, Cr(III), and the lanthanide, which suppresses QTM.\(^2,3^3\) These Cr–Dy SMMs, 11(Me), 11(Et), 11(Bu) and 12, display hysteresis of the magnetization vs. applied field at temperatures as high as 2.2. K.

This core topology can also be found in the family [Ln\(_n\)M\(_n\)(OH)\(_2\)Cym\(_n\)COO\(_3\)](THF)\(_3\)] (Ln = Dy 13, Ho; Cym = \((\mu_5-, C_\text{H}_\text{$_5$})\text{C}_\text{H}_\text{$_3$})\text{Mn}$$\text{(CO)}_\text{$_6$}\)) where only the dysprosium analogue presents SMM properties.\(^3^4\) In 2011 Powell and co-workers reported the Type B defective dicubane complexes [Ni\(_2\)Ln\(_2\)(L)\(_2\)NO\(_3\)\(_2\)] (Ln = Dy 14(S), Tb; L = 2-(2-hydroxy-3-methoxy-benzylidene-amino)phenol; S = MeOH and DMP)\(^3^5\) where the two dysprosium complexes, 14(MeOH) and 14(DMF) are SMMs and the blocking temperature seems to be modulated by the coordination environment around the Ni(n) ions. The structure of these complexes is somewhat different from the SMMs reported by Murray and co-workers; now the dysprosium ions are not part of the central [M\(_2\)O\(_2\)] unit, but at the tips of the molecule, resulting in two fairly well separated ions, nearly magnetically independent. In 2012 the same group reported [Co\(_2\)Dy\(_2\)L\(_4\)](NO\(_3\)\(_2\)THF)\(_2\)] (15),\(^3^6\) a Type B defective dicubane, which was also a SMM. Cobalt is in the oxidation state Co(n), and is paramagnetic. The authors show how single-ion blocking of the Dy(iii) ions occurs at higher temperatures with a crossover to molecular exchanged-based blocking at low temperatures. For 15 there are two differentiated thermally-activated regimes with effective barriers of 82 and 11 K. Hysteresis loops were clearly observed up to 3 K. In this very interesting paper the authors unambiguously assign the large energy barrier to the relaxation of the Dy(iii) ions and the low temperature behaviour to the exchange-blocked relaxation where the 3d–4f coupling dominates.

A number of higher nuclearity 3d–4f complexes have been reported, many of them SMMs with diverse structures and ligands. Most of these have relatively small energy barriers. An exception is an enneanuclear complex [Mn\(_2\)Dy\(_n\)O\(_3\)(mdea)\(_2\)–(mdeaH)\(_2\)Piv\(_n\)](NO\(_3\))\(_2\)H\(_2\)O\(_2\)] (16(Dy)) reported by Powell et al.; the complex possesses an energy barrier of 38.6 K and displays magnetisation vs. field hysteresis loops up to 1.9 K.\(^3^7\) The Tb(iii), Ho(iii) and Y(iii) (16(Tb), 16(Ho) and 16(Y)) analogues are also SMMs, but with smaller energy barriers. The diamagnetic
Y(III) analogue highlights the fact that the Mn₅ unit also contributes to the SMM behaviour. The complex \([\text{Mn}_{21}\text{DyO}_{20}(\text{OH})_2(\text{BuCOO})_{17}(\text{HCOO})_4(\text{NO}_3)_4(\text{H}_2\text{O})_7]\) (17) shown in Fig. 4, reported by Christou and co-workers in 2011, also shows hysteresis of the magnetization up to 3 K and has a large energy barrier of 74 K.²⁸ In complex 17 a single Dy(III) is at one vertex of a \([\text{Mn}_4\text{O}_8]\) cubane, which is surrounded by a Mn–O core that is probably a strong contributor to the SMM properties. The 3d–4f SMM containing the largest number of Dy(III) ions is complex 18 with ten Dy(III) and two cobalt ions, \([\text{Dy}_{10}\text{Co}_2(\text{L})_4(\text{MeCOO})_{16}(\text{SCN})_2(\text{MeCN})_2(\text{H}_2\text{O})_2]\) from 2011 by Tang and his group.²⁹ Complex 18 contains ten dysprosium ions and two Co(n) in a wheel-like arrangement. Below 10 K the authors report slow relaxation with a crossover from single-ion relaxation to the exchanged coupled system due to the Dy(n) centres. Two complexes have been reported with eight dysprosium ions. \([\text{H}_2\text{O}]\text{[Cu}_{12}\text{Dy}_4(\text{Ph}_3\text{CPO}_3)_{16}(\text{Ph}_3\text{CPO}_3\text{H})_{12}(\text{MeCOO})_{12}(\text{MeCOOH})_{12}(\text{OH})_{12}(\text{NO}_3)_2(\text{H}_2\text{O})_2]\) was reported in 2010 by Winpenny and co-workers.³⁰ The Gd(III) analogue is not an SMM, thus the anisotropy is provided by the dysprosium ions. As usual the exchanged coupled system due to the Dy(III) centres. Two more examples of 3d–4f SMM complexes fail to provide the best magnetic properties. In 2015 Dehnen and co-workers reported the octanuclear complex \([\text{Mn}_9\text{O}_9(\text{saO})_6(\text{MeO})_6\text{Bo}(\text{MeOH})_6(\text{H}_2\text{O})_2]\) (21(Tb)) with an energy barrier of 103 K.⁴⁵ This complex displayed a high blocking temperature of 3.1 K. The lanthanum analogue, 21(La) was also an SMM with a large blocking temperature and relatively large energy barrier of 32 K. Brehin and co-workers reported the same year complex \([\text{Mn}_{21}\text{DyO}_{20}(\text{OH})_2(\text{OAc})_4(\text{MeCOO})_{20}(\text{HCOO})_4(\text{NO}_3)_4(\text{H}_2\text{O})_7]\) (22) with a similar core, but different ligands and an energy barrier of 24 K.⁴⁶ The energy barrier was also high for the terbium complex 22(Tb), \(U_{\text{eff}} = 46 K\). In this case clearly the anisotropy of the terbium ion has actually boosted the SMM properties of the \([\text{Mn}_n\text{Dy}_1\text{O}_5]\) unit, showing how the right combination of anisotropic 3d metal and anisotropic lanthanide can lead to better, improved SMMs. The high nuclearity, high symmetry (D₂) complex 23(Tb) reported by Tong and co-workers, \([\text{Mn}_{21}\text{DyO}_{20}(\text{OH})_2(\text{OAc})_4(\text{mpea})_{14}(\text{MeO})_6\text{Bo}(\text{MeOH})_6(\text{H}_2\text{O})_2]\) has lower symmetry (C₁) and a smaller energy barrier value of 19.6 K. In this particular case, the dysprosium analogue 23(Dy) has a smaller anisotropy barrier. Complex 25, \([\text{Mn}_7\text{DyO}_6(\text{mdea})_6(\text{mdeaH})_6(\text{MeO})_6(\text{H}_2\text{O})_2]\), with a large blocking temperature and relatively large energy barrier of 3.1 K. The lanthanum analogue highlights the fact that the Mn₅ unit also contributes to the SMM behaviour.

There are several 3d–Tb SMMs synthesized with Cu, the most relevant is complex 27, \([\text{Cu}_{12}\text{Dy}_{21}(\text{OH})_2(\text{H}_2\text{O})(\text{NO}_3)_2]\)–\([\text{ClO}_4]_2\), with \(U_{\text{eff}} = 25 K\), a tetragonal pyramid with a large and flexible ligand bis(carboxyethyl)isocyanurate.⁴⁹ On the other hand, we can see an example of a family of trinuclear \([\text{Cu}_{12}\text{Tb}_{21}(\text{OH})_2(\text{H}_2\text{O})(\text{NO}_3)_2]\)–\([\text{ClO}_4]_2\) complexes like \([\text{Cu}_{12}\text{Dy}_{21}(\text{OH})_2(\text{H}_2\text{O})(\text{NO}_3)_2]\)–\([\text{ClO}_4]_2\) with different assemblies, showing values of anisotropy barrier between 18 K and 23 K.²⁰ A one-dimensional chain of units like that of complexes 28 and 29 behaves as a single-chain magnet (SCM).²⁰ Two more examples of Cu–Tb SMMs are complex 30, \([\text{Cu}_{12}\text{Tb}_{21}(\text{OH})_2(\text{H}_2\text{O})(\text{NO}_3)_2]\)–\([\text{ClO}_4]_2\) or complex 31, \([\text{Cu}_{12}\text{Tb}_{21}(\text{OH})_2(\text{H}_2\text{O})(\text{NO}_3)_2]\)–\([\text{ClO}_4]_2\) with an hexameric macrocycle ligand.⁵³ Complex 32 is a rare example...
of a heterotrimetallic coordination complex of formula \([\text{Cu(H}_2\text{L})(\text{MeO})\text{H}]_2\text{Th(H}_2\text{O})_6\text{DMF}]_{0.43}\text{Fe(CN)}_6\), exhibiting an energy barrier of 13 K\(^5^4\). Some other SMMs reported in the reference list are an octanuclear complex \([\text{Cu(Tb)}_2\text{Cl}_4\text{O}_4\text{MeO}_3\text{H}]_{1.5}\text{NO}_3\) complex 33, described as an oblate wheel that has an energy barrier of 15.6 K\(^5^5\) or complex 34, a dinuclear Cu–Tb complex, \([\text{LCu(OMe)Tb(thd)}_2\text{]}_2\), with a Schiff base ligand with \(U_{\text{eff}} = 13\) K\(^5^6\). The complex \([\text{Tb}_2\text{Ni}_4\text{Cl}_2\text{O}_4\text{OH}_2\text{MeO}_2\text{MeOH}_6\text{Cl}_2\text{Cl}_4\text{O}_4\text{h}_3\), 35(Tb), is a defective dicubane complex with Schiff base ligands. Complex 35(Tb) exhibits an anisotropy barrier of 30 K. The magnetic data obtained for the Ni–Y[m] analogue 35(Y) demonstrated that the 4f metal contribution to the SMM properties was dominant\(^5^7\). Complex 36 is another rare heterotrinuclear 3d–4f–5d complex containing tungsten, \([\text{[LMe}_2\text{Ni(}2\text{H}_2\text{O})\text{Tb(thd)}_2\text{]}_2\text{]}_2\text{]}_2\text{]}_2\), showing an anisotropy barrier of 15 K. It is worth pointing out that none of the reported 3d–Tb(m) SMMs report magnetization vs. field hysteresis loops above 2.0 K.

**Sm, Ho, Er and Yb 3d–4f SMMs**

Even though dysprosium and terbium are undoubtedly the two lanthanide ions that have provided better SMMs, there are some interesting examples with other lanthanide ions. Samarium, with a less than half-filled shell and the smallest 4f SMM, \([\text{Fe}_{12}\text{Sm}_{4}\text{O}_{10}\text{(OH)}_4\text{(PhCO}_2\text{)}_2\text{]}_4\), \(35\text{(Sm)}\), is a defective dicubane complex with Schiff base ligands. Complex 35(Sm) exhibits an anisotropy barrier of 30 K. The magnetic data obtained for the Ni–Y[m] analogue 35(Y) demonstrated that the 4f metal contribution to the SMM properties was dominant. Complex 36 is another rare heterotrinuclear 3d–4f–5d complex containing tungsten, \([\text{[LMe}_2\text{Ni(}2\text{H}_2\text{O})\text{Tb(thd)}_2\text{]}_2\text{]}_2\text{]}_2\), showing an anisotropy barrier of 15 K. It is worth pointing out that none of the reported 3d–Tb(m) SMMs report magnetization vs. field hysteresis loops above 2.0 K.

**Challenges ahead and concluding remarks**

The biggest challenge still remains, to raise the blocking temperature of new SMMs, no matter whether we talk about transition metals, 4f or 3d–4f SMMs. We have included here some complexes that are claimed as SMMs but for which no maxima in the out-of-phase ac magnetic susceptibility or hysteresis of the magnetization vs. field are reported. Clearly, the limited access to experiments at temperatures below 2 K is an obstacle in this respect. However, we expect that when new 3d–4f SMMs with higher blocking temperatures are reported, this fact will cease to be a problem and there will be less ambiguity as to the physical properties of reported species: both ac magnetic susceptibility out-of-phase maxima and magnetization vs. field hysteresis should be observed to claim a complex is an SMM. With lanthanide SMMs the effective energy barriers have been greatly increased, up to hundreds of Kelvin in several mononuclear SMMs, Tb–phthalocyanine derivatives, pure or doped in diamagnetic yttrium matrices, but this has not been accompanied by a real increase in blocking temperatures, thus hampering the potential application of lanthanide SMMs. This problem might be overcome by 3d–4f SMMs. Several examples report large effective energy barriers (at least of the order of those reported for 3d SMMs) that are in a few cases accompanied by relatively high blocking temperatures, such as those reported by Murray and co-workers for \([\text{Cr}_{2}\text{Dy}_{2}\text{]}\) SMMs, with \(T_b = 2.2\) K, and particularly relevant, the \([\text{Mn}_3\text{Tb}_2\text{]}\) reported by Dehnen and co-workers with \(U_{\text{eff}} = 103\) K and \(T_b = 3.5\) K\(^3^3,3^4\). The \(\text{Cr(III)}–\text{Dy(III)}\) significant magnetic interaction is claimed to be the key factor in quenching QTM and it is directly related to the anisotropy barrier, thus opening up a challenging new route to control SMM properties of \([\text{Cr(III)}–\text{Dy(III)}\)] ions. Could this be exploited for other 3d–4f SMMs? It is still a big synthetic challenge to prepare 3d–4f complexes with strong magnetic coupling between the 3d and 4f ions, but this might be a great goal to have in mind. The synthetic methods clearly offer a rich variety of products, with different levels of control in the design of the prepared complexes. There is not a clear picture of preferred ligands to prepare 3d–4f SMMs, and complexes are reported with all kinds of ligands, but polycalix ligands and Schiff bases of salicylaldehyde appear in many of the reported complexes.
The advances in the theoretical understanding of the magnetic properties of the lanthanide ions and their 3d–4f complexes are still lagging behind the advances in the synthesis of new complexes. There is still a lot to learn about heterometallic 3d–4f complexes, especially about the magnetic coupling between 3d and 4f metals. We strongly believe the study of 3d–4f interactions as it becomes more common, even in dinuclear model complexes, will provide good ideas for the design of new 3d–4f SMMs. Furthermore, two main trends of design of new 3d–4f SMMs have emerged as the most plausible to provide better 3d–4f SMMs in the near future: isolated lanthanide ions with a 3d metallo ligand, as in the [Mn$_2$Dy]$^+$$^-$ reported by Christou and co-workers, with $T_b = 3.0$ K; or 3d–4f complexes with strong magnetic coupling between the metals to suppress QTM. Also a combination of these approaches emerges as a good option: a 3d SMM with strong coupling between a unique lanthanide ion and a 3d metallo ligand with large S that would help in quenching the QTM, thus increasing the blocking temperature.

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