Study of the Production of Radioactive Isotopes through Cosmic Muon Spallation in KamLAND

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Radioactive isotopes produced through cosmic muon spallation are a background for rare event detection in ν detectors, double-beta-decay experiments, and dark-matter searches. Understanding the nature of cosmogenic backgrounds is particularly important for future experiments aiming to determine the pep and CNO solar neutrino fluxes, for which the background is dominated by the spallation production of ¹³C. Data from the Kamioka Liquid scintillator Anti-Neutrino Detector (KamLAND) provides valuable information for better understanding these backgrounds, especially in liquid scintillator, and for checking estimates from current simulations based upon MUSIC, FLUKA, and Geant4. Using the time correlation between detected muons and neutrton captures, the neutron production yield in the KamLAND liquid scintillator is measured to be \((2.8 ± 0.3) \times 10^{-4} \, n/(\mu \cdot (g/cm^2))\). For other isotopes, the production yield is determined from the observed time correlation related to known isotope lifetimes. We find some yields are inconsistent with extrapolations based on an accelerator muon beam experiment.

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I. INTRODUCTION

Cosmic ray muons and their spallation products are potential sources of background for neutrino detectors, double-beta decay experiments, and dark-matter searches, even when the detectors are deployed underground. Characterizing cosmic-ray-muon-induced backgrounds, particularly the secondary neutrons and radioactive isotopes produced by muon-initiated spallation processes, is essential for interpreting these experiments.

Liquid-scintillator detectors such as KamLAND, Borexino [1–3], CANDLES IV [4, 5], SNO+ [6, 7], LENs [8], and LENA [9] are designed to detect low-energy phenomena. In organic liquid scintillator (LS), energetic muons and subsequent showers interact mostly with $^{12}$C, the most abundant nucleus heavier than $^1$H in the LS, generating neutrons and isotopes by electromagnetic or hadronic processes. The muon-initiated spallation of carbon targets is a matter of primary interest.

Isotope production by muon-initiated spallation has been studied by an earlier experiment [10] using the CERN Super Proton Synchrotron (SPS) muon beam. The energy dependence was studied with 100 and 190 GeV incident muons. The production yield at other energies is estimated from this data by extrapolation, assuming a power-law dependence on the muon energy. Direct measurements of the production yield by underground detectors such as LSD [11], LVD [12], and Borexino [13] were compared to calculations exploiting simulations based on MUSIC [14], FLUKA [15, 16], and Geant4 [17, 18]. Particular attention was paid to neutron production since isotope production measurements are difficult with the small scintillator masses used in these detectors. KamLAND, owing to its larger mass $\sim$1 kton of LS – does not suffer from this difficulty and is well placed to study a variety of isotopes of interest.

This paper presents the neutron and isotope production rates in KamLAND from muon-initiated spallation based upon data collected from 5 March 2002 to 12 May 2007. The results are compared to simulations and other experiments. These comparisons provide important information for validating Monte Carlo simulations.

II. DETECTOR DESCRIPTION AND PERFORMANCE

KamLAND is located under the peak of Ikenoyama (Ike Mountain, 36.42°N, 137.31°E), and the vertical rock overburden is approximately 2700 meters water equivalent (m.w.e.). A schematic diagram of KamLAND is shown in Fig. 1. KamLAND consists of an active detector region of approximately 1 kton of ultra-pure LS contained in a 13-m-diameter spherical stainless-steel outer vessel. A buffer comprising of 57% isoparaffin and 43% dodecane oils by volume fills the region between the balloon and the surrounding 18-m-diameter spherical stainless-steel outer vessel to shield the LS from external radiation. The specific gravity of the buffer oil (BO) is adjusted to be 0.04% lower than that of the LS. An array of photomultiplier tubes (PMTs), 1325 specially developed fast PMTs masked to 17-inch diameter and 554 older 20-inch diameter PMTs reused from the Kamiokande experiment [21], are mounted on the inner surface of the outer containment vessel, providing 34% photocathode coverage. During the period from 5 March 2002 to 27 February 2003 the photo-cathode coverage was only 22%, since the 20-inch PMTs were not operated. A 3 mm thick
acrylic barrier at 16.6-m-diameter helps prevent radon emanating from the PMT glass from entering the BO. The inner detector (ID), consisting of the LS and BO regions, is surrounded by a 3.2 kton water Čerenkov detector instrumented with 225 20-inch PMTs. This outer detector (OD) absorbs γ-rays and neutrons from the surrounding rock and enables tagging of cosmic-ray muons.

The KamLAND front-end electronics (FEE) system is based on the Analog Transient Waveform Digitizer (ATWD) [22] which captures PMT signals in 128 10-bit digital samples at intervals of 1.5 ns. Each ATWD captures three gain levels of a PMT signal to obtain a dynamic range from one photoelectron (p.e.) to 1000 p.e. Each ATWD takes 27 µs to read out, so two are attached to each PMT channel to reduce dead time. The FEE system contains discriminators set at 0.15 p.e. (∼0.3 mV) threshold which send a 125 ns long logic signal to the trigger electronics. The trigger electronics counts the number of ID and OD PMTs above the discriminator threshold with a sampling rate of 40 MHz and initiates readout when the number of 17-inch ID PMTs above the discriminator threshold (N₁₇) exceeds the number corresponding to ∼0.8 MeV deposited energy. The trigger system also issues independent readout commands when the number of OD PMTs above threshold exceeds a preset number.

The energy can be estimated from the Nmax parameter, defined as the maximum value of N₁₇ in a 200 ns period following the trigger command. However, the offline analysis takes full advantage of the information stored in the digitized PMT signals by identifying individual PMT pulses in the waveform information that is read out. The time and integrated area (called charge) are computed from the individual pulses. For each PMT, the average charge corresponding to a single p.e. is determined from single-pulse waveforms observed in low occupancy events. The ID PMT timing is calibrated with light pulses from a dye laser (∼1.2 ns pulse width), injected at the center of the detector through an optical fiber. The vertices of spatially localized low-energy (<30 MeV) events are estimated by comparing calculated time-of-flight of optical photons from the hypothetical vertex to the measured arrival times at the PMTs in KamLAND.

The reconstructed energies of events are calibrated with γ sources: 203Hg, 68Ge, 65Zn, and 60Co; and with n + γ sources: 241Am+9Be and 210Po+13C [23]. These are deployed at various positions along the vertical axis of the detector and occasionally off the vertical axis within 5.5 m from the detector center [24]. Such calibrations cover energies between 0.28 and 6.1 MeV. The energy calibration is aided with studies of background contaminants 40K and 208Tl, 213Bi, 212Po and 214Bi, 214Po sequential decays, 12B and 12N spallation products, and γ’s from thermal neutron captures on 1H and 12C.

The visible energy (Evis) of an event is computed from the measured light yield. Specifically, Evis is the number of detected p.e. after corrections for PMT variation, dark noise, solid angle, shadowing by suspension ropes, optical transparency, and scattering properties in the LS. The relationship between Evis and the deposited energy (Edep) of γ’s, e±’s, protons, and α’s is non-linear and modeled as a combination of Birks-quenched scintillation [25, 26] and Čerenkov radiation. The scale is adjusted so that Evis is equal to Edep for the 2.225 MeV γ-ray from neutron capture on 1H. The observed energy resolution is ∼7.4%/√Evis(MeV) for the period without the 20-inch PMTs, and ∼6.5%/√Evis(MeV) for the rest of the data.

The calibration sources are also used to determine systematic deviations in position reconstruction by comparison with the source’s known position. This comparison gives an average position reconstruction uncertainty of less than 3 cm for events with energies in the range 0.28 to 6.1 MeV.

### III. COSMIC RAY MUONS

A digital map [27] of the topological profile of Ikenoyama is shown in Fig. 2. The vertical overburden at KamLAND is approximately 1000 meters of rock and the minimum overburden corresponding to a nearby valley is approximately 900 meters.

Cosmic ray muons are identified either by the large amount of scintillation and Čerenkov light detected by the ID PMTs, or by the Čerenkov light detected by the OD PMTs. Muons...
crossing the ID (ID muons) are selected by requiring that one of the following conditions was satisfied:

1. \( \mathcal{L}_{\text{ID}} \geq 10000 \text{ p.e.} \) \((\sim 30 \text{ MeV})\)
2. \( \mathcal{L}_{\text{ID}} \geq 500 \text{ p.e.} \) and \( N_{\text{OD}} \geq 5 \)

where \( \mathcal{L}_{\text{ID}} \) is the total light yield measured by the ID 17-inch PMTs, and \( N_{\text{OD}} \) is the number of OD PMTs with signals above threshold. Approximately 93% of the ID muons satisfy the first selection criterion. The ID muon track is reconstructed from arrival times of the first-arriving Čerenkov or scintillation photons at the PMTs. Since for relativistic muons the wavefront of the scintillation light proceeds at the Čerenkov angle, and since muons generate enough light to generate photoelectrons in every PMT, by restricting the fit to the first-arriving photons both Čerenkov and scintillation photons can be treated identically. The observed muon track is then established by minimizing time-of-flight deviations from hypothetical muon tracks. The fit converges for 97% of all ID muon events. The majority of the events that are not reconstructed are believed to be multiple muons, or muons accompanied by large electromagnetic or hadronic showers for which the tracking model is not valid.

Muons passing only through the BO produce mostly Čerenkov light, whereas muons passing through the LS generate both Čerenkov and scintillation light. Figure 3 shows the correlation between the light yield \( \mathcal{L}_{\text{ID}} \) and the shortest distance between the reconstructed muon track and the center of KamLAND (impact parameter). The boundary at 650 cm between the BO and LS regions is evident. The correlations between \( \mathcal{L}_{\text{ID}} \) and the reconstructed muon track length in the BO and LS regions \( L_{\text{BO}} \) and \( L_{\text{LS}} \), respectively) are plotted in Figs. 4a and 4b. A linear trend, corresponding to minimum ionizing muons, is apparent in both distributions. The slope of each line is the light yield per unit length in the respective material. In BO, where the light is predominantly Čerenkov, the light yield per unit length is found to be \( \langle d\mathcal{L}_{\text{C}}/dX \rangle = 31 \pm 2 \text{ p.e./cm} \); the fit was restricted to path lengths above 700 cm since fits at shorter path lengths are complicated by the presence of PMTs which may obstruct some of the emitted light. In the LS we obtain

\[
\frac{d\mathcal{L}_{\text{S}}}{dX} = \frac{\mathcal{L}_{\text{ID}} - L_{\text{BO}} \langle d\mathcal{L}_{\text{C}}/dX \rangle}{L_{\text{LS}}} = 629 \pm 47 \text{ p.e./cm}, \quad (1)
\]

where \( d\mathcal{L}_{\text{S}}/dX \) includes the Čerenkov light created in the LS. The muons in Fig. 4 generating light yields above the baseline linear trend are likely to involve secondary particles. We define an excess light yield parameter \( \Delta\mathcal{L} \),

\[
\Delta\mathcal{L} = \mathcal{L}_{\text{ID}} - L_{\text{BO}} \langle d\mathcal{L}_{\text{C}}/dX \rangle - L_{\text{LS}} \langle d\mathcal{L}_{\text{S}}/dX \rangle \quad (2)
\]

for the purpose of describing showering muons associated with secondary particles.

The LS muon rate is estimated by selecting muons with \( \mathcal{L}_{\text{ID}} > 4 \times 10^4 \text{ p.e.} \) and impact parameter \(< 650 \text{ cm}\). The light yield cut has a negligible inefficiency for LS muons, while the impact parameter cut eliminates LS muons that reconstruct outside the balloon because of the resolution of the fitter algorithm. This provides the lower limit on the LS muon rate. An upper limit is established by removing the impact parameter cut and increasing the \( \mathcal{L}_{\text{ID}} \) cut to \( > 10^5 \text{ p.e.} \) to eliminate muons that pass through the BO without transversing the LS. This cut again has a small inefficiency for LS muons but does not eliminate muons which
Ref. [33] made a detailed determination of KamLAND, and SuperKamiokande detectors. The authors of there are previous estimates of the mean energy (initiated spallation is energy dependent. KamLAND does not contributions were taken into account.

flux from Ref. [29] and cross sections from Ref. [30] gives

\[ \sim 10^{5} \text{p.e.}, \] which is equivalent to a \( \sim 3 \text{ GeV} \) threshold, the

rate of showering muons in the LS is \( \sim 0.03 \text{ Hz} \). It is possible that some atmospheric neutrino interactions leak into the LS muon sample. However, an estimate with the neutrino flux from Ref. [29] and cross sections from Ref. [30] gives less than \( 4 \times 10^{-5} \text{ Hz} \) from atmospheric neutrinos. The muon track length distribution is shown in Fig. 5. The measured average track length is \( L_\mu = 878 \text{ cm} \), in agreement with the calculated value of \( L_\mu = 874 \pm 13 \text{ cm} \) where the non-spherical corrections to the balloon shape and the muon angular distributions were taken into account.

The production yield of radioactive isotopes from muon-initiated spallation is energy dependent. KamLAND does not measure the muon energy, so it is estimated from simulation. There are previous estimates of the mean energy (\( \bar{E}_\mu \)) ranging from 198 GeV [31] to 285 GeV [10, 32] at the Kamiokande, KamLAND, and SuperKamiokande detectors. The authors of Ref. [33] made a detailed determination of \( J_\mu \) and \( \bar{E}_\mu \) using the MUon SImulation Code (MUSIC) [14] to transport muons, generated according to the Modified Gaisser Parameterization [34] sea-level muon flux distribution, through a digital profile of Ikenoyama. Although this calculation reproduces the zenith and azimuthal angular distributions observed by KamLAND, as shown in Fig. 10 of Ref. [33], it overestimates \( J_\mu \) by 14% relative to this work. The calculation assumed a homogeneous rock entirely of the Inishi type (see Table II of Ref. [33] for the chemical composition), but other rock types, such as granite, limestone, and several types of metamorphic rock, are common in Ikenoyama in unknown quantities [35]. Table II shows the result of calculations of \( J_\mu \) and \( \bar{E}_\mu \) using the same simulation method of Ref. [33] but for Standard Rock [36, 37] and Generic Skarn, a generic mixture of rock types found in a skarn-type mine like that of Ikenoyama. Here, Generic Skarn is defined to be 70% Granite and 30% Calcite by weight. The \( J_\mu \) values range from \( 4.90 \text{ m}^{-2} \text{h}^{-1} \), for \( 2.75 \text{ g/cm}^3 \) specific gravity Generic Skarn, to \( 6.71 \text{ m}^{-2} \text{h}^{-1} \), for \( 2.65 \text{ g/cm}^3 \) specific gravity Inishi Rock, while \( \bar{E}_\mu \) varies from 254 GeV to 268 GeV. The value of \( J_\mu \) for \( 2.70 \text{ g/cm}^3 \) specific gravity Standard Rock is \( 5.38 \text{ m}^{-2} \text{h}^{-1} \), in excellent agreement with our measured value. The value of \( \bar{E}_\mu \) for this rock is 259 GeV. We take \( \bar{E}_\mu = 260 \pm 8 \text{ GeV} \), where the uncertainty is chosen to cover the full range for the various rock types.

IV. SPALLATION NEUTRON YIELD

Most of the neutrons produced in the KamLAND LS capture on hydrogen or carbon atoms. The capture cross section varies inversely with respect to velocity, and the mean neutron capture time (\( \tau_n \)) is constant with respect to energy. The capture time (\( t \)) distribution is exponential, \( P(t) \propto e^{-t/\tau_n} \). A calculation using the elemental composition of the KamLAND LS shown in Table I and the thermal neutron capture cross sections from Ref. [38] gives \( \tau_n = 206 \mu \text{s} \). This calculation indicates that 99.5% of the neutrons capture on \(^1\text{H} \), while the remainder capture mostly on \(^{13}\text{C} \). The probability for capture on the other isotopes in the LS, such as \(^{13}\text{C} \), is \( 2 \times 10^{-4} \) or less.

Neutrons produced by muon-initiated spallation in the LS can be identified by the characteristic capture \( \gamma \)-rays. Figure 6 shows the \( E_{\text{vis}} \) distributions in signal (150 \( \leq \Delta T < 1000 \mu \text{s} \)) and background (4150 \( \leq \Delta T < 5000 \mu \text{s} \)) coincidence windows following some muons, where \( \Delta T \equiv (t - t_\mu) \) is the time elapsed since the muon’s passage. The \( E_{\text{vis}} \) distribution clearly shows peaks from neutron captures on \(^1\text{H} \) (2.225 MeV) and \(^{12}\text{C} \) (4.9 MeV), which are not evident in the

| Rock Type                  | \( J_\mu \) (m\(^{-2}\)h\(^{-1}\)) | \( \bar{E}_\mu \) (GeV) |
|----------------------------|---------------------------------|------------------------|
| Inishi Rock                | 5.66 to 6.71                    | 262 to 268             |
| Standard Rock              | 4.95 to 5.83                    | 256 to 262             |
| Generic Skarn              | 4.90 to 5.82                    | 254 to 260             |
| This Measurement           | 5.37 \( \pm \) 0.41             | —                      |
background. Figure 7 shows the $\Delta T$ distribution for events within the LS volume and with $1.8 < E_{vis} < 2.6$ MeV, which includes the single $2.225$ MeV $\gamma$-ray emitted by neutron capture on $^1H$. For $\Delta T < 1000\mu s$, there is a clear deviation from the exponential distribution due to the overload that large muon signals produce on individual electronics channels and to the dead time in the system arising from the very high event multiplicity following the muon. Both effects intervene in events that are quite different to those which KamLAND was designed to record.

The number of neutrons produced by muon-initiated spallation in the LS is established by a binned maximum likelihood fit [39] to the data in Fig. 7 using the function

$$r(t) = \frac{N_n}{\tau_n} e^{-(t-t_n)/\tau_n} + r_B,$$

where $N_n$ is the total number of neutron captures associated with the selected events, $\tau_n$ is the mean neutron capture time, and $r_B$ is the background rate, which is assumed to be approximately constant in the region of interest ($\Delta T < 2500\mu s$) due to the low muon rate ($\sim 0.2$ Hz).

To avoid the electronics-induced distortions, the fit is restricted to the region $\Delta T > 1300\mu s$. The parameters $N_n$ and $r_B$ are free, but the mean capture time is constrained to $\tau_n = 207.5 \pm 2.8\mu s$ with a gaussian penalty function. This mean capture time is determined from two independent measurements of $\tau_n$: a $^{241}$Am+$^9$Be calibration source, and an analysis of a sample of neutrons generated by clean muons. These clean muons are identified by a multiplicity parameter $\eta_T$ defined to be the number of trigger commands that follow the muon within a 10 ms period. Fits to the $\Delta T$ distribution of the subset of neutrons selected with various limits on $\eta_T$ demonstrate that muons with $\eta_T < 10$ give unbiased fit residuals. The value of $\tau_n$ is $207.5 \pm 0.3\mu s$ from the clean muon sample and $205.2 \pm 0.5\mu s$ from the $^{241}$Am+$^9$Be source data. The observed $2.3\mu s$ discrepancy between these values is not completely understood, but is suspected to be caused by neutrons from the $^{241}$Am+$^9$Be source that capture on the stainless-steel source capsule. In this analysis, we use the value $\tau_n = 207.5\mu s$ from the non-showering muon sample with an uncertainty of $\pm 2.8\mu s$ covering both measurements. The fit shown in Fig. 7 in the region $\Delta T \geq 1300\mu s$ results in $N_n = (4.2 \pm 0.3) \times 10^6$ and $\chi^2$/d.o.f.$= 98/118$.

The actual number of neutrons $N_n$ produced by muon-initiated spallation is related to the fit result $N_n$ by an efficiency $\epsilon_n$,

$$N_n = \frac{N_n}{\epsilon_n},$$

that accounts for other neutron-eliminating nuclear reactions like $^{12}$C$(n,p)$ and neutron losses at the LS-BO boundary. This efficiency is calculated using the MUSIC-based muon simulation described in Sec. III and the Geant4-based Monte Carlo of KamLAND described in Sec. VI A. The muon simulation generates muons with a 3-momentum distribution appropriate to the KamLAND site inside Ikenoyama. Muon

| Effect                              | Value        |
|-------------------------------------|--------------|
| Neutron Eliminating Reactions (e.g. $(n, p)$) | 96.3 ± 3.7% |
| Neutron Captures on $^1H$           | 99.5 ± 0.1% |
| LS-BO Boundary                      | 93.3 ± 2.0% |
| Electronics Dead Time Effects       | > 98%        |
| Combined Efficiency                 | 89.4 ± 3.8% |

TABLE III: Summary of the dominant contributions to the neutron detection efficiency.
transport is then modeled through KamLAND. In this simulation neutrons are created and destroyed; neutrons that survive to thermalization are tracked until they are captured. Energy scale nonlinearities and the finite γ-ray resolution are included. The effects of detector boundaries are taken into account. The component of the efficiency related to the electronics dead time for high-multiplicity events is measured by comparing the number of recorded waveforms with the number of trigger commands. The correction is measured to be less than 2% for ∆T ≥ 1300 μs. In counting neutrons, we use the definition by which the (n, 2n) reaction generates one new neutron and the (n, n′) reaction generates no new neutrons. This method gives an efficiency εn = 89.4 ± 3.8% (broken down in Table III).

Using Eqn. 4 we find \( N_n = (4.7 ± 0.4) \times 10^6 \), which is then used to extract the neutron production yield:

\[
Y_n = \frac{N_n}{R_\mu T_L \rho L_\mu}, \tag{5}
\]

where \( R_\mu = 0.198 ± 0.014 \text{Hz} \) is the measured rate of LS muons, \( T_L = 1.24 \times 10^8 \text{s} \) is the detector live time, \( \rho = 0.780 ± 0.001 \text{g/cm}^3 \) is the density of the LS, and \( L_\mu = 874 ± 13 \text{cm} \) is the calculated mean muon track length. The resulting neutron production yield is \( Y_n = (2.8 ± 0.3) \times 10^{-7} \text{n}/(μ \cdot \text{g/cm}^2) \), with 64 ± 5% of the neutrons produced by events classified as showering muons.

V. SPALLATION ISOTOPE YIELD

The method for determining the yields of spallation-generated isotopes is similar to the neutron analysis described in Sec. IV. Spallation-generated isotopes are identified by their decay time relative to their creation, and by their decay energy. The decay time, \( ∆T = t - t_\text{vis} \), is calculated for each event relative to all previous muons. Usually, several different isotopes decay in a given time window. The number of each isotope produced, \( N_i \), is obtained from a binned maximum likelihood fit [39] to the ∆T distribution, using the function:

\[
r(t) = \sum_i N_i \tau_i e^{-(t-t_\text{vis})/\tau_i} + r_B, \tag{6}
\]

where \( N_i \) is related to \( N_i \) by an event selection efficiency (\( N_i = N_i/ε_i \)), \( τ_i \) is the mean lifetime, and \( r_B \) is a constant rate of uncorrelated background events that are in accidental coincidence with muons.

In terms of \( N_i \), the spallation production yield for isotope \( i \) is equal to

\[
Y_i = \frac{N_i}{R_\mu T_L \rho L_\mu}, \tag{7}
\]

where \( R_\mu, T_L, \rho, \) and \( L_\mu \) are defined as in Eqn. 5. The spallation production rate for isotope \( i \) is equal to

\[
R_i = \frac{N_i}{ρ V_T T_L}, \tag{8}
\]

where \( V_T = 1171 ± 25 \text{m}^3 \) is the target volume.

![FIG. 8: Background-subtracted \( E_\text{vis} \) spectrum above 4 MeV. Signal and background events are taken from 2 ≤ ∆T < 60 ms and 502 ≤ ∆T < 560 ms, respectively. The production rate of \(^{12}\text{B}\) (τ = 29.1 ms, \( Q = 13.4 \text{MeV} \)) is estimated to be 54.8 ± 1.5 (kt-on-day)−1, from fitting Eqn. 6 to the ∆T distribution shown in the inset. This fit has a χ²/d.o.f. = 509/495. The production rate \(^{12}\text{N}\) (τ = 15.9 ms, \( Q = 17.3 \text{MeV} \)) is estimated to be 2.2 ± 0.5 (kt-on-day)−1, from a similar fit to the ∆T distribution for events with 14 ≤ \( E_\text{vis} \) < 20 MeV, where the higher \( E_\text{vis} \) threshold is imposed to exclude \(^{12}\text{B}\).]

A. \(^{12}\text{B} \) and \(^{12}\text{N} \)

\(^{12}\text{B}\) (τ = 29.1 ms, \( Q = 13.4 \text{MeV} \)) [40] β− decay and \(^{12}\text{N}\) (τ = 15.9 ms, \( Q = 17.3 \text{MeV} \)) [40] β+ decay candidate events are selected via cuts on \( E_\text{vis} \) and ∆T. The inset in Fig. 8 shows the distribution of ∆T for events with 4 ≤ \( E_\text{vis} \) < 20 MeV and 2 ≤ ∆T < 500 ms. A binned maximum likelihood fit to the ∆T distribution using Eqn. 6, and including the long-lived isotopes \(^8\text{He} \), \(^9\text{Li} \), \(^9\text{C} \), and \(^{12}\text{N} \) as contaminants, yields \( N(\text{^{12}B}) = (5.94 ± 0.10) \times 10^4 \) events with χ²/d.o.f. = 509/495. Longer lived isotopes give roughly constant decay rates on this time scale and fit out as a component of \( r_B \). A similar fit to the ∆T distribution for events with 14 ≤ \( E_\text{vis} \) < 20 MeV, where the higher \( E_\text{vis} \) threshold is imposed to exclude \(^{12}\text{B} \), gives \( N(\text{^{12}N}) = (2.8 ± 0.3) \times 10^2 \) events. A comparison with the predicted \( E_\text{vis} \) spectrum is shown in Fig. 8. The \( E_\text{vis} \) spectra are predicted from the allowed \(^{12}\text{B} \) and \(^{12}\text{N} \) β decay spectra, taking into account the KamLAND detector response, normalized to the observed \( N(\text{^{12}B}) \) and \( N(\text{^{12}N}) \). The detector response model includes the energy nonlinearities and boundary effects described in Sec. II.

The inefficiency in identifying \(^{12}\text{B} \) and \(^{12}\text{N} \) candidates is dominated by the \( E_\text{vis} \) cut. The efficiencies calculated by integrating the predicted \( E_\text{vis} \) spectra over the selection window give ε(\(^{12}\text{B} \)) = 82.9 ± 0.7% and ε(\(^{12}\text{N} \)) = 9.3 ± 1.6%, where the errors come from the uncertainty in the detector response. Using Eqn. 7, the resultant isotope production yields are calculated to be \( Y(\text{^{12}B}) = (42.9 ± 3.3) \times 10^{-7} (μ \cdot \text{g/cm}^2) \)−1 and
The fit to the $\Delta T$ distribution has a $\chi^2$/d.o.f. = 95/91.

The efficiency is given by the product of $\epsilon_E$ for the $E_{\text{vis}}$ selection and $\epsilon_S$ for the muon-$^8$Li/$^8$B spatial correlation. By integrating the expected spectra over $4 \leq E_{\text{vis}} < 20$ MeV, $\epsilon_E(^8\text{Li})$ and $\epsilon_E(^8\text{B})$ are estimated to be $77.6 \pm 0.9\%$ and $88.4 \pm 0.7\%$, respectively. For showering muons, $\epsilon_S$ is 100% because no correlation requirement is imposed. For non-showering muons, $\epsilon_S$ is estimated from the $^{12}\text{B}/^{12}\text{N}$ analysis (Sec. V A), and the systematic error arising from variations between isotopes is estimated with a FLUKA simulation. Figure 10 shows the impact parameter ($\Delta L$) distribution for the $^{12}\text{B}/^{12}\text{N}$ candidate events for non-showering muons ($\Delta L < 10^6$ p.e.). We find that $91.6 \pm 4.3\%$ of the candidates are within 3 m of the muon track. This fraction is the value of $\epsilon_S(^{12}\text{B})$ for $\Delta L < 3$ m.

In order to obtain $\epsilon_S(^8\text{Li})$ and $\epsilon_S(^8\text{B})$, an additional correction for the difference between the muon-$^8\text{Li}/^8\text{B}$ and the muon-$^{12}\text{B}$ spatial correlations is applied. This correction is derived from the FLUKA simulation described in Sec. VI B. The simulation does not include the uncertainties in the muon track and the isotope decay vertex reconstruction; it is only used to study the isotope dependence of $\epsilon_S$. The range of values of $\epsilon_S$ from FLUKA for different spallation isotopes is used to estimate the systematic error that should be added to $\epsilon_S(^{12}\text{B})$ in order to obtain a common $\epsilon_S$ for all spallation isotopes. The resulting value, $\epsilon_S = 91.6 \pm 8.4\%$, is used for estimating the $^8\text{Li}$ and $^8\text{B}$ (and, later the $^6\text{C}$, $^8\text{He}$, and $^9\text{Li}$) yields.

Combining the above analyses, we obtain $Y(^{6}\text{Li}) = (12.2 \pm 2.6 \times 10^{-7} (\mu \cdot (g/cm^2))^{-1}$ and $Y(^{8}\text{B}) = (8.4 \pm 2.4 \times 10^{-7} (\mu \cdot (g/cm^2))^{-1}$. The isotope production rates are $R(^{6}\text{Li}) = 15.6 \pm 3.2$ (kton-day)$^{-1}$ and $R(^{8}\text{B}) = 10.7 \pm 2.9$ (kton-day)$^{-1}$. The contour plots in Fig. 11 show the correlation between $^8\text{Li}$ and $^8\text{B}$ due to their similar lifetimes. $^8\text{Li}$ and $^8\text{B}$ are identified primarily by their energy spectra.
Since a reduced volume is used in the signal window from a signal in the $0.5 \times 0.5 \times 0.5$ cm$^3$ Neutron detection requirement, which is calculated with the Geant4-based Monte Carlo simulation described in Sec. VI A. The resulting efficiencies $\epsilon(^{8}\text{He}) = 14.9 \pm 1.0\%$ and $\epsilon(^{8}\text{Li}) = 46.1 \pm 1.1\%$ include the appropriate $\beta^{-}n$ branching fractions. Since a reduced volume is used in this analysis, the 1.6% fiducial volume uncertainty from Ref. [44] is included in the above efficiencies. The resulting yields are $Y(^{8}\text{He}) = (0.7 \pm 0.4) \times 10^{-7} \ (\mu \cdot (g/cm^2))^{-1}$ and $Y(^{8}\text{Li}) = (2.2 \pm 0.2) \times 10^{-7} \ (\mu \cdot (g/cm^2))^{-1}$. The production rates are $R(^{8}\text{He}) = 1.0 \pm 0.5 \ (kton \cdot day)^{-1}$ and $R(^{8}\text{Li}) = 2.8 \pm 0.2 \ (kton \cdot day)^{-1}$. The contour plots in Fig. 13 show the correlation between $^{8}\text{Li}$ and $^{8}\text{He}$. The procedure for calculating the $^{8}\text{He}$ and $^{9}\text{Li}$ selection efficiency is the same as for the $^{8}\text{Li}^{8}\text{B}$ selection efficiency analysis (Sec. V B), except for the correction for the neutron detection requirement, which is calculated with the Geant4-based Monte Carlo simulation described in Sec. VI A. The neutron is excited states of $^{8}\text{He}$ with a $16 \pm 1\%$ branching ratio [41], and $^{9}\text{Li}$ decays to neutron-unstable excited states of $^{8}\text{Be}$ with a $50.8 \pm 0.9\%$ branching ratio [41]. The neutron is identified by the 2.225 MeV $\gamma$-ray from radiative capture on $^{1}\text{H} \ (1.8 \leq E_{vis} \leq 2.6 \text{MeV})$. The $\gamma$-ray is required to be within 200 cm and 1.0 ms of the $^{8}\text{He}^{8}\text{B}^{8}\text{Li}$ $\beta^{-}$ decay candidate. Finally, the $^{8}\text{He}^{8}\text{Li}$ analysis is performed using a 5.5-m-radius spherical fiducial volume to reduce the number of accidental coincidences between the $\beta^{-}$ decay candidate and external $\gamma$-ray backgrounds near the balloon.

The inset in Fig. 12 shows the $\Delta T$ distribution for the events that satisfy the criteria outlined above. Figure 12 also shows the residual $E_{vis}$ distribution corresponding to the subtraction of a background spectrum in the 5.002 $\leq \Delta T < 6\text{ s}$ window from a signal in the 0.002 $\leq \Delta T < 1\text{ s}$ window. The expected $E_{vis}$ distributions for $^{8}\text{He}$ and $^{9}\text{Li}$ are calculated by incorporating the KamLAND response and adjusting for the energy deposited by the thermalizing neutron from $^{8}\text{He}$ or $^{9}\text{Li}$ decay. $N(^{8}\text{He})$ and $N(^{9}\text{Li})$ are determined from a simultaneous maximum likelihood fit to the $\Delta T$ distribution and a chi-square fit to the $E_{vis}$ distribution. For the fit to the $E_{vis}$ distribution, the uncertainty in the energy scale parameters are treated in the same manner as the $^{8}\text{Li}^{8}\text{B}$ analysis described in Sec. V B.

C. $^{8}\text{He}$ and $^{9}\text{Li}$

$^{8}\text{He}$ ($\tau = 171.7 \text{ ms}, \ Q = 10.7 \text{ MeV}$) [41] and $^{9}\text{Li}$ ($\tau = 257.2 \text{ ms}, \ Q = 13.6 \text{ MeV}$) [41] $\beta^{-}$ decay candidate events are selected according to the cuts $1 \leq E_{vis} \leq 13 \text{ MeV}$ and $\Delta T < 10\text{ s}$, and by the detection of a neutron following the $\beta^{-}$ decay event. $^{8}\text{He}$ decays to neutron-unstable excited states of $^{8}\text{Li}$ with a $16 \pm 1\%$ branching ratio [41], and $^{9}\text{Li}$ decays to neutron-unstable excited states of $^{8}\text{Be}$ with a $50.8 \pm 0.9\%$ branching ratio [41]. The neutron is identified by the 2.225 MeV $\gamma$-ray from radiative capture on $^{1}\text{H} \ (1.8 \leq E_{vis} \leq 2.6 \text{MeV})$. The $\gamma$-ray is required to be within 200 cm and 1.0 ms of the $^{8}\text{He}^{8}\text{B}^{8}\text{Li}$ $\beta^{-}$ decay candidate. Finally, the $^{8}\text{He}^{8}\text{Li}$ analysis is performed using a 5.5-m-radius spherical fiducial volume to reduce the number of accidental coincidences between the $\beta^{-}$ decay candidate and external $\gamma$-ray backgrounds near the balloon.

The inset in Fig. 12 shows the $\Delta T$ distribution for the events that satisfy the criteria outlined above. Figure 12 also shows the residual $E_{vis}$ distribution corresponding to the subtraction of a background spectrum in the 5.002 $\leq \Delta T < 6\text{ s}$ window from a signal in the 0.002 $\leq \Delta T < 1\text{ s}$ window. The expected $E_{vis}$ distributions for $^{8}\text{He}$ and $^{9}\text{Li}$ are calculated by incorporating the KamLAND response and adjusting for the energy deposited by the thermalizing neutron from $^{8}\text{He}$ or $^{9}\text{Li}$ decay. $N(^{8}\text{He})$ and $N(^{9}\text{Li})$ are determined from a simultaneous maximum likelihood fit to the $\Delta T$ distribution and a chi-square fit to the $E_{vis}$ distribution. For the fit to the $E_{vis}$ distribution, the uncertainty in the energy scale parameters are treated in the same manner as the $^{8}\text{Li}^{8}\text{B}$ analysis described in Sec. V B.

The procedure for calculating the $^{8}\text{He}$ and $^{9}\text{Li}$ selection efficiency is the same as for the $^{8}\text{Li}^{8}\text{B}$ selection efficiency analysis (Sec. V B), except for the correction for the neutron detection requirement, which is calculated with the Geant4-based Monte Carlo simulation described in Sec. VI A. The resultant efficiencies $\epsilon(^{8}\text{He}) = 14.9 \pm 1.0\%$ and $\epsilon(^{9}\text{Li}) = 46.1 \pm 1.1\%$ include the appropriate $\beta^{-}n$ branching fractions. Since a reduced volume is used in this analysis, the 1.6% fiducial volume uncertainty from Ref. [44] is included in the above efficiencies. The resulting yields are $Y(^{8}\text{He}) = (0.7 \pm 0.4) \times 10^{-7} \ (\mu \cdot (g/cm^2))^{-1}$ and $Y(^{8}\text{Li}) = (2.2 \pm 0.2) \times 10^{-7} \ (\mu \cdot (g/cm^2))^{-1}$. The production rates are $R(^{8}\text{He}) = 1.0 \pm 0.5 \ (kton \cdot day)^{-1}$ and $R(^{8}\text{Li}) = 2.8 \pm 0.2 \ (kton \cdot day)^{-1}$. The contour plots in Fig. 13 show the correlation between $^{8}\text{Li}$ and $^{8}\text{He}$.
D. 9C

The inset in Fig. 14 shows the $\Delta T$ distribution for all events with visible energy $12 \leq E_{vis} < 20$ MeV. The analysis region ($0.2 \leq \Delta T < 0.6$ s) contains events from $^9$C ($\tau = 182.5$ ms, $Q = 16.5$ MeV) [41] $\beta^+$ decay. $N(^9$C) is determined from a simultaneous binned maximum likelihood fit to the $\Delta T$ distribution and a chi-square fit to the $E_{vis}$ distribution. The uncertainty in the energy scale parameters are treated in the same manner as the $^8$Li/$^8$B analysis described in Sec. V B. In this fit, $^8$Li, $^8$B, and $^8$Be are treated as possible contaminants, the amounts are constrained to the values obtained in the previously described analyses. This constraint includes the correlation between $^8$Li and $^8$B shown Fig. 11. By integrating the theoretical $^9$C $E_{vis}$ spectrum, we obtain the efficiency for the $12 \leq E_{vis} < 20$ MeV cut of $\epsilon(^9$C) = 7.2 ± 1.0%. Combining this with the above results gives $Y(^9$C) = $(3.0 \pm 1.2) \times 10^{-7} (\mu \cdot \text{g/cm}^2)^{-1}$ and $R(^9$C) = $3.8 \pm 1.5$ (kton-day)$^{-1}$.

E. $^{11}$C

The production of $^{11}$C ($\beta^+$-decay, $\tau = 29.4$ min, $Q = 1.98$ MeV) [40] through muon-initiated spallation is usually accompanied by a neutron, allowing identification by the triple coincidence of the primary muon, the spallation neutron, and the subsequent $\beta^+$ [13, 45]. The $^{11}$C $\beta^+$ decays are selected in the range $1.4 \leq E_{vis} < 2.0$ MeV and are preceded by a detected muon that is accompanied by at least one neutron capture, identified by the 2.225 MeV $\gamma$-ray from capture on $^3$H. The $\gamma$-ray is required to be in the time window

$10 \leq \Delta T < 2500 \mu$s relative to the muon. In order to reduce the background, a 7-m-diameter fiducial volume is used. To avoid inefficiencies from run boundaries and the long lifetime of $^{11}$C, the first 5 hours of the typically 24-hour long run are not used in the selection of the $^{11}$C candidates. The number of muon-$^{11}$C coincidences is extracted from the $\Delta T$ distribution for all events that meet the criteria, shown in the inset of Fig. 15.

The efficiency determination takes into account the visible energy range for $^{11}$C $\beta^+$-decay, $22.7 \pm 3.6\%$, and the previously discussed neutron detection efficiency (Sec. IV). The efficiency also takes into account a correction for $^{11}$C production modes, designated invisible modes, which do not produce neutrons [13, 45]. To measure this correction, muonic $^{11}$C event pairs were selected with and without the neutron requirement for a subset of the data where the $^{11}$C candidate is required to be within 50 cm of the muon track; restricting the study to a subset of the data mitigated the reduced signal-to-background ratio associated with relaxing the neutron requirement. The number of muon-$^{11}$C coincidences in each case was extracted from a fit of Eqn. 6 to the corresponding $\Delta T$ distribution. The visible mode efficiency, $\epsilon_{vis}$, taken as the ratio of the number of muon-$^{11}$C pairs with one or more neutrons to the number of muon-$^{11}$C pairs without the neutron requirement, is $88.4 \pm 2.4\%$. Applying the correction for post muon electronics effects, the visible mode fraction is $96.3 \pm 2.0\%$, consistent with Ref. [45], which obtains $\epsilon_{vis} = 95.6\%$ for 285 GeV muons. Due to the relatively long $^{11}$C lifetime, we have also considered the effect of diffusion. An analysis of $^{222}$Rn that was accidentally introduced into the center of KamLAND during the deployment of a calibration device shows that
efficient is $89\%$. The fit to the $\Delta T$ distribution has a $(\chi^2$/d.o.f.$= 80/56$).

The production rate of $^{10}$C ($\tau = 27.8$ s, $Q = 3.65$ MeV) is $21.1 \pm 1.8$ (kton-day)$^{-1}$, as determined by fitting Eqn. 6 to the $\Delta T$ distribution shown in the inset. The fit to the $\Delta T$ distribution has a $(\chi^2$/d.o.f.$= 80/56$).

the diffusion speed is approximately 1 mm/h. From this study the effect of $^{11}$C diffusion on the efficiency is estimated to be less than 0.5%. Combining this with the above results gives $Y(^{11}\text{C}) = (866 \pm 153) \times 10^{-7} (\mu \cdot (\text{g/cm}^2))^{-1}$ and $R(^{11}\text{C}) = 1106 \pm 178 \text{ (kton-day)}^{-1}$.

**F. $^{10}$C**

As with $^{11}$C, the production of $^{10}$C($\beta^+$-decay, $\tau = 27.8$ s, $Q = 3.65$ MeV) [41] through muon-initiated spallation is usually accompanied by a neutron, so the selection criterion requiring a triple coincidence of the primary muon, the neutron, and $^{10}$C candidate is used. The neutron is identified by the 2.225 MeV $n + ^{1}\text{H}$ capture $\gamma$-ray. The number of $^{10}$C candidates is determined from fitting Eqn. 6 to the $\Delta T$ distribution for all events identified as $^{10}$C, shown in the inset in Fig. 16. $^{11}$Be is a potential background for this $^{10}$C analysis, but the correction to the $^{10}$C yield is estimated to be less than 1% to due to the low $^{11}$Be production rate and the neutron coincidence requirement. The efficiency for the visible energy cut $2.0 \leq E_{vis} < 4.0$ MeV is 73.5 $\pm$ 3.2%. The visible mode efficiency is $\varepsilon_{vis} = 90.7 \pm 5.5%$ after correcting for the electronics effects following muons. The final efficiency is $89.6 \pm 5.5%$. The resulting isotope yield is $Y(^{10}\text{C}) = (16.5 \pm 1.9) \times 10^{-7} (\mu \cdot (\text{g/cm}^2))^{-1}$ and the production rate is $R(^{10}\text{C}) = 21.1 \pm 1.8 \text{ (kton-day)}^{-1}$.

**G. $^{11}$Be**

The $^{11}$Be $\beta^-$ decay ($\tau = 19.9$ s, $Q = 11.5$ MeV) [40] events are selected according to $5.5 \leq E_{vis} < 16.0$ MeV. The $E_{vis}$ cut efficiency is 63.4%. The inset in Fig. 17 shows the $\Delta T$ distributions for the events only after showering muons. For non-showering muons, a tighter muon track cut $\Delta L < 1$ m is applied in order to reduce the background rate. The track cut efficiency is estimated from the $^{12}$B candidates using an analysis similar to Sec. VB. The resulting isotope yield is $Y(^{11}\text{Be}) = (1.1 \pm 0.2) \times 10^{-7} (\mu \cdot (\text{g/cm}^2))^{-1}$ and production rate is $R(^{11}\text{Be}) = 1.4 \pm 0.3 \text{ (kton-day)}^{-1}$.

**VI. MONTE CARLO SIMULATIONS**

The Geant4 and FLUKA simulations are used to reproduce the measurements from KamLAND. While Geant4 is used only to simulate neutron production, both neutron and light isotope production are tested with FLUKA.

**A. Geant4**

Geant4 is a widely used toolkit for performing particle tracking simulations on an event-by-event basis. A description of the available physics processes included is given in Refs. [17] and [18]. Here we compare the Geant4 (version 9.1) prediction for neutron yield by spallation with the results obtained in Sec. IV. We use the physics list QGS_BIC, developed by the Geant4 group to support the binary cascade (BIC) model at lower energies (below 10 GeV for $p$ and $n$, and below 1.2 GeV for $\pi$). This treatment is also appropriate for

**FIG. 16:** Background-subtracted $E_{vis}$ spectrum above 2 MeV, where signal and background are selected using $10 \leq \Delta T < 90$ s and $190 \leq \Delta T < 270$ s, respectively. The production rate of $^{10}$C ($\tau = 27.8$ s, $Q = 3.65$ MeV) is $21.1 \pm 1.8$ (kton-day)$^{-1}$, as determined by fitting Eqn. 6 to the $\Delta T$ distribution shown in the inset. The fit to the $\Delta T$ distribution has a $(\chi^2$/d.o.f.$= 80/56$).

**FIG. 17:** Background-subtracted $E_{vis}$ spectrum above 5.5 MeV for showering muons ($\Delta C > 10^6$ p.e.), where signal and background are selected by $8 \leq \Delta T < 60$ s and $408 \leq \Delta T < 460$ s, respectively. The production rate of $^{11}$Be ($\tau = 19.9$ s, $Q = 11.5$ MeV) for showering muons is estimated to be $1.0 \pm 0.2$ (kton-day)$^{-1}$ by fitting Eqn. 6 to the $\Delta T$ distribution shown in the inset $(\chi^2$/d.o.f.$= 37/41$).
the simulation of interactions of nucleons and ions. At higher
energies, a quark-gluon string (QGS) model is applied for the
hadronic interactions. Neutron elastic and inelastic interac-
tions below 20 MeV are described by a high-precision data-
driven model (NeutronHP). The G4EmExtraPhysics
physics list is also used to model the photo-nuclear and muon-
nuclear interaction processes, which dominate the neutron
production by muons in the simulation.

To estimate the neutron production yield as a function of
muon energy in Geant4, mono-energetic muons of several
energies are injected at the center of a generic hydrocarbon
block of thickness 40 m. The region more than 10 m away
from the edges of the block is analyzed to avoid boundary ef-
fects. As shown in Fig. 20, the neutron production yields pre-
dicted by Geant4 are systematically lower than experiment,
with the exception of point (F) from LVD [12]. These results
are consistent with previous work [46].

A Monte Carlo simulation based upon Geant4 and
MUSIC (described in Sec. III) is used to study neutrons pro-
duced by muon-induced spallation in the material outside of
the KamLAND ID. Some of these neutrons have sufficient en-
ergy to enter the ID where they thermalize and capture. They
can be identified by the coincidence of a prompt signal (for
example from n + p elastic scattering) and a delayed signal
from the capture γ-ray. This is the same inverse beta decay
reaction signature used for νe detection, \( ν_e + p → e^+ + n \),
where the \( e^+ \) is the prompt signal and the γ-ray from neu-
tron capture is the delayed signal and therefore a potential
background. These neutrons are also a background for dark
matter experiments that employ nuclear recoils as a detection
method.

The primary purpose of this Monte Carlo simulation is
to estimate the rate of untagged fast neutrons, i.e. neutrons
produced by muons where the muon is undetected by the
KamLAND ID. A few of these muons are detected by the
OD, either from the Čerenkov radiation produced by the muon
itself (tracked muons), or by accompanying electromagnetic

\[ \frac{\text{Events/m}^3}{\text{cm}^2} \]

FIG. 18: \( E_{\text{vis}} \) distribution of the prompt events from candidates
identified as neutrons produced by muon-induced spallation in the
material outside of the KamLAND ID.

and hadronic showers that enter the OD (untracked muons).
A prompt signal in coincidence with these tracked and un-
tracked muons followed by a delayed capture gamma signal
identifies candidates for untagged fast neutrons.

Fast neutrons generated by untracked muons are produced
primarily in the rock surrounding the OD, whereas fast neu-
trons from tracked muons are primarily produced in the
water of the OD. It is shown in the Monte Carlo that the tracked
and untracked muons give distinguishable OD visible energy dis-
tributions. A comparison between Monte Carlo and measure-
ment of the distribution of the number of OD PMTs with sig-
als above threshold for tracked and untracked muons reveals
a deficiency of fast neutrons from untracked muons, consistent
with the underproduction of neutrons by Geant4 in concrete
reported in Ref. [47].

Figure 18 shows the \( E_{\text{vis}} \) distribution of the prompt events
from the Monte Carlo simulation of tracked and untracked
muons compared to the measured data. The measured data
and the Monte Carlo simulation corresponds to an equal live
time exposure of 1368 days. Figure 19 shows the radial distri-
bution of the prompt events relative to the KamLAND center.
Both the measured data and the Geant4 simulation exhibit an
exponential attenuation of the neutrons as they penetrate fur-
ther into the detector, with the simulation yielding an attenua-
tion length of \( 69 \pm 2 \text{ g/cm}^2 \), consistent with the measurement
\( 70 \pm 2 \text{ g/cm}^2 \).

B. FLUKA

FLUKA is a mature code that models nuclear and partic-
le physics processes from thermal neutrons to heavy ion
collisions [15, 16]. It has been used previously to model
TABLE IV: Simulation of neutron and light isotope production in the KamLAND LS by muon-initiated spallation with FLUKA: the isotope production yield by mono-energetic (260 GeV) $\mu^-$; the ratio of the production yields by mono-energetic $\mu^+$ compared to $\mu^-$; the ratio of the production yields by a $\mu^-$ spectrum that matches Fig. 6 (KamLAND curve) in Ref. [33] compared to mono-energetic $\mu^-$; the power law exponent for the production yield $Y(E_{\mu}) \propto E_{\mu}^{\alpha}$ from a fit to the yields from mono-energetic $\mu^-$ with $10 \leq E_{\mu} \leq 350$ GeV; and the primary process for producing the isotope. The uncertainties are statistical.

| Simulated Production Yield | Ratio of Simulated Production Yields for | Power Law Exponent | Primary Process |
|----------------------------|----------------------------------------|-------------------|----------------|
| $(\times 10^{-7} \mu \cdot (g/cm^2))^{-1}$ | $\mu^+/\mu^-$ | Spectrum / Mono-Energetic | $\pi^{-}+H,^{12}C$ |
| $n$ | 2344 ± 4 | 0.969 ± 0.002 | 0.912 ± 0.003 | 0.779 ± 0.001 | $\pi^{-}+H,^{12}C$ |
| $^{11}C$ | 460.8 ± 1.7 | 0.971 ± 0.005 | 0.913 ± 0.006 | 0.703 ± 0.002 | $^{12}C(\gamma, n)$ |
| $^{7}Be$ | 116.8 ± 0.9 | 0.986 ± 0.011 | 0.945 ± 0.011 | 0.684 ± 0.004 | $^{12}C(\gamma, \alpha)$ |
| $^{9}Be$ | 44.63 ± 0.53 | 0.960 ± 0.018 | 0.891 ± 0.019 | 0.825 ± 0.007 | $^{12}C(n, ^3He)$ |
| $^{12}B$ | 30.85 ± 0.44 | 0.970 ± 0.021 | 0.936 ± 0.022 | 0.828 ± 0.009 | $^{12}C(n, p)$ |
| $^{8}Li$ | 23.42 ± 0.39 | 0.927 ± 0.026 | 0.936 ± 0.025 | 0.821 ± 0.010 | $^{12}C(n, p\alpha)$ |
| $^{10}C$ | 21.13 ± 0.37 | 0.982 ± 0.025 | 0.915 ± 0.027 | 0.810 ± 0.010 | $^{12}C(\pi^+, np)$ |
| $^{6}He$ | 13.40 ± 0.29 | 0.916 ± 0.035 | 0.918 ± 0.035 | 0.818 ± 0.013 | $^{12}C(n, 2p^3He)$ |
| $^{8}B$ | 6.40 ± 0.20 | 0.996 ± 0.045 | 0.915 ± 0.050 | 0.804 ± 0.019 | $^{12}C(\pi^+, ^3H, ^4H)$ |
| $^{9}Li$ | 3.51 ± 0.15 | 0.856 ± 0.074 | 0.842 ± 0.078 | 0.801 ± 0.026 | $^{12}C(\pi^+, ^3He)$ |
| $^{9}C$ | 1.49 ± 0.10 | 0.850 ± 0.114 | 0.949 ± 0.102 | 0.772 ± 0.039 | $^{12}C(\pi^+, ^3H)$ |
| $^{12}N$ | 0.86 ± 0.07 | 0.963 ± 0.128 | 1.006 ± 0.120 | 0.921 ± 0.045 | $^{12}C(p, n)$ |
| $^{11}Be$ | 0.94 ± 0.08 | 0.842 ± 0.145 | 0.804 ± 0.161 | 0.753 ± 0.051 | $^{12}C(n, 2p)$ |
| $^{8}He$ | 0.35 ± 0.05 | 0.964 ± 0.200 | 0.576 ± 0.372 | 0.926 ± 0.078 | $^{12}C(\pi^+, ^3H)$ |
| $^{13}B$ | 0.31 ± 0.04 | 1.020 ± 0.197 | 1.062 ± 0.176 | 0.742 ± 0.075 | $^{12}C(n, p)$ |
| $^{15}O$ | 0.05 ± 0.02 | 1.250 ± 0.379 | 1.635 ± 0.234 | 0.793 ± 0.244 | $^{10}O(\gamma, n)$ |
| $^{13}N$ | 0.06 ± 0.02 | 1.500 ± 0.272 | 1.190 ± 0.401 | 1.120 ± 0.220 | $^{13}C(p, n)$ |

Muon-initiated spallation in liquid scintillator [31, 46, 52, 53]. We use FLUKA Version 2006.3b to model neutron and light isotope production from muon-initiated spallation in KamLAND. A 40 m radius by 40 m high cylinder of KamLAND liquid scintillator is used in the simulation; the concentric inner cylinder of 20 m radius and 20 m length is used for analysis.

To estimate neutron production yield as a function of muon energy in FLUKA, mono-energetic beams of $\mu^-$ ranging from 10 GeV to 350 GeV were simulated as in Ref. [31, 46, 52, 53]. Care is taken not to double count neutrons involved in reactions like $(n, 2n)$. The results of this simulation are included in Fig. 20. The neutron production yield of this FLUKA simulation is 10% lower than previous work by [31, 46, 52, 53], but the power law dependence on muon energy ($E_{\mu}^{\alpha}$, where $\alpha = 0.77$) is consistent. Different scintillator compositions were studied, but could not explain the deficit. This deficit is insignificant compared to the discrepancy between these simulations and the data.

The production of light isotopes was studied in the same simulation. The results, including the primary production process and power law exponent, are summarized in Table IV. For some isotopes the primary production process is much larger than any secondary processes, as in the case of $^{12}B$. For other isotopes the primary production process is only slightly larger than the secondary processes, as is the case for $^9Li$. The isotopes produced primarily by $\gamma$ interactions, $^{11}C$ and $^{10}C$, show the weakest dependence on muon energy. In comparison, $^{12}N$ and $^{13}N$, where the primary production mechanism is by $p$ interactions, show the strongest dependence on muon energy.

The use of a mono-energetic $\mu^-$ beam overestimates the production of neutrons and light isotopes. Simulations using a mono-energetic $\mu^+$ beam and a beam with the energy spectrum from Ref. [33] were also run. The simulations show that the production yield for $\mu^+$ relative $\mu^-$ is on average 0.96±0.01 for the light isotopes. This reduction is expected since $\mu^-$ may capture creating spallation products while $\mu^+$ may not. This ratio, combined with the $\mu^+$ to $\mu^-$ ratio at KamLAND, leads to a correction to the flux of 0.98±0.06 for light isotopes and 0.98±0.005 for neutrons. The reduced production yield due to averaging over the muon spectrum is on average 0.92±0.02 for the light isotopes, which is slightly higher than the correction factor suggested by Ref. [10]. The results of the FLUKA simulations for KamLAND presented in Table V include these two corrections.

VII. DISCUSSION

The isotope production yields from muon-initiated spallation in a liquid scintillator target were investigated at CERN by earlier experiment [10] using the SPS muon beam with muon energies of 100 GeV and 190 GeV. Based on those cross section measurements and the predicted muon energy spectrum at the KamLAND site, we calculated the isotope pro-
TABLE V: Summary of the neutron and isotope production yields from muon-initiated spallation in KamLAND. The results of the FLUKA calculation shown in this table include corrections for the muon spectrum and the $\mu^+ / \mu^-$ composition of the cosmic-ray muon flux.

| Lifetime in KamLAND LS | Radiation Energy | Spallation Production Yield ($\times 10^{-7} (\mu \cdot (g/cm^2))^{-1}$) | FLUKA calc. | Fraction from showering $\mu$ |
|------------------------|------------------|-------------------------------------------------|--------------|-----------------------------|
| $\eta$                 | 207.5 $\mu$s     | 2.225 MeV (capt. $\gamma$)                       | ---          | ---                         |
| $^{12}$B               | 29.1 ms          | 13.4 MeV ($\beta^-$)                             | ---          | 2097 ± 13                   |
| $^{12}$N               | 15.9 ms          | 17.3 MeV ($\beta^+$)                             | ---          | 2787 ± 311                  |
| $^8$Li                 | 1.21 s           | 16.0 MeV ($\beta^- \alpha$)                      | ---          | 64 ± 5%                     |
| $^8$B                  | 1.11 s           | 18.0 MeV ($\beta^+ \alpha$)                      | ---          | 42.9 ± 3.3                  |
| $^9$C                  | 182.5 ms         | 16.5 MeV ($\beta^+$)                             | ---          | 68 ± 2%                     |
| $^9$He                 | 171.7 ms         | 10.7 MeV ($\beta^- \gamma n$)                    | 1.0 ± 0.3    | 77 ± 14%                    |
| $^9$Li                 | 257.2 ms         | 13.6 MeV ($\beta^- \gamma n$)                    | 3.16 ± 0.25  | 77 ± 6%                     |
| $^{11}$C               | 29.4 min         | 1.98 MeV ($\beta^+$)                             | 421 ± 68     | 62 ± 10%                    |
| $^{10}$C               | 27.8 s           | 3.65 MeV ($\beta^- \gamma$)                      | 54 ± 12      | 76 ± 6%                     |
| $^{11}$Be              | 19.9 s           | 11.5 MeV ($\beta^-$)                             | < 1.1        | 74 ± 12%                    |
| $^{8}$He               | 1.16 s           | 3.51 MeV ($\beta^-$)                             | 7.5 ± 1.5    | ---                         |
| $^7$Be                 | 76.9 day         | 0.478 MeV (EC $\gamma$)                          | 107 ± 21     | ---                         |

The production yields for the isotopes from muon-initiated spallation in KamLAND are provided in Table V. On average, the yields from the showering muons (≈15% of all muons), whose excess light yield parameter ($\Delta L$, Eqn. 2) is greater than $10^6$ p.e. (≈3 GeV), constitute 70 ± 2% of the yield from all muons. The production yield for $^{11}$C is the largest, and its measured yield is larger than the FLUKA calculation by a factor of ≈2. The Borexino collaboration also reported a similar discrepancy [2, 3], that is consistent with what is observed in KamLAND. Some measured production yields, such as $^8$Li and $^{10}$C, deviate significantly from estimates based on the muon beam experiment, indicating that perhaps estimation by extrapolation is not sufficient. All isotope yields are consistent within an order of magnitude.

VIII. SUMMARY

We have analyzed KamLAND data to measure production yields of radioactive isotopes and neutrons through muon-initiated spallation in liquid scintillator. The neutron production yield is evaluated to be $(2.8 \pm 0.3) \times 10^{-4} n/(\mu \cdot (g/cm^2))$, which is higher than the expectation from Monte Carlo simulations based on Geant4 and FLUKA. Some isotope production yields are found to be inconsistent with extrapolations – based on a power law dependence with respect to muon energy – of results from muon beam experiments.

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