Z’ and Dark Matter

Yann Mambrini
Laboratoire de Physique Théorique, Université Paris-Sud, F-91405 Orsay, France
E-mail: yann.mambrini@th.u-psud.fr

Abstract. Several constructions motivate the existence of a dark \( U(1)_D \) gauge boson which interacts with the Standard Model only through its kinetic mixing or loop induced processes. We describe two typical examples with specific signatures in particular we show that a region with relatively light WIMPS, \( M_{Z_D} < 40 \) GeV and a kinetic mixing \( 10^{-4} < \delta < 10^{-3} \) is not yet excluded by the last experimental data and seems to give promising signals in a near future. We also show that conditions from anomaly cancelation generate tri-vector couplings \( Z_D Z\gamma \) leading to a specific gamma ray line observable by FERMI telescope.

1. Introduction
Neutral gauge sectors with an additional dark \( U(1)_D \) symmetry in addition to the Standard Model (SM) hypercharge \( U(1)_Y \) and an associated \( Z_D \) are among the best motivated extensions of the SM, and give the possibility that a dark matter candidate lies within this new gauge sector of the theory [1]. The new vector boson \( Z_D \) can interact with the SM, even if no SM fermions are directly charged under the additional gauge symmetry. This interaction can occur via mixed kinetic terms between the SM’s hypercharge field strength and the new abelian field strength [2, 3, 4, 5] or through couplings generated by counter-term to preserve the anomaly cancelation condition [6, 7]. Whereas the former couplings can give significant signals in direct detection experiment even fitting the last DAMA [8] or COGENT [9] excesses [10, 11, 12, 13] or INTEGRAL 511 keV line [14], the latter can give rise to a gamma-ray line observable in satellite telescopes [15, 16, 17, 18]. A summary of gamma-ray line constraints can be found in [19], and within a supersymmetric [20] in extra-dimensional framework in [21]. A light \( Z’ \) can be justified in leptophobic constructions [22] and even be an interpretation of the \( Wjj \) anomaly observed by CDF [23]. Using LEP/Tevatron constraints, a higher dimensional approach can be found in [25] whereas a summary of all the constraints has been studied in [26]

2. The dark kinetic mixing
The matter content of any \( dark\ U(1)_D \) extension of the SM can be decomposed into three families of particles:

- The \textit{Visible sector} is made of particles which are charged under the SM gauge group \( SU(3) \times SU(2) \times U(1)_Y \) but not charged under \( U(1)_D \) (hence the \textit{dark} denomination for this gauge group)
- the \textit{Dark sector} is composed by the particles charged under \( U(1)_D \) but neutral with respect of the SM gauge symmetries. The dark matter (\( \psi_0 \) candidate is the lightest particle of the \textit{dark sector}
• The Hybrid sector contains states with SM and $U(1)_D$ quantum numbers. These states are fundamental because they act as a portal between the two previous sector through the kinetic mixing they induce at loop order.

From these considerations, it is easy to build the effective lagrangian generated at one loop:

$$\mathcal{L} = \mathcal{L}_\text{SM} - \frac{1}{4} \tilde{B}_{\mu\nu} \tilde{B}^{\mu\nu} - \frac{1}{4} \tilde{X}_{\mu\nu} \tilde{X}^{\mu\nu} - \frac{\delta}{2} \tilde{B}_{\mu\nu} \tilde{X}^{\mu\nu} + i \sum_i \psi_i \gamma^\mu D_\mu \psi_i + i \sum_j \Psi_j \gamma^\mu D_\mu \Psi_j$$

(1)

$B_\mu$ being the gauge field for the hypercharge, $X_\mu$ the gauge field of $U(1)_D$ and $\psi_i$ the particles from the hidden sector, $\Psi_j$ the particles from the hybrid sector, $D_\mu = \partial_\mu - i(\bar{q} \gamma^\gamma \tilde{B}_\mu + q_D \tilde{g}_D \tilde{X}_\mu + g T^a W_\mu^a)$, $T^a$ being the $SU(2)$ generators, and

$$\delta = \frac{\bar{q} \gamma^\gamma \tilde{B}_\mu + q_D \tilde{g}_D \tilde{X}_\mu + g T^a W_\mu^a}{16\pi^2} \sum_j q^j \bar{q}^j \log \left( \frac{m_j^2}{M_j^2} \right)$$

(2)

with $m_j$ and $M_j$ being hybrid mass states [4]. Notice that the sum is on all the hybrid states, as they are the only ones which can contribute to the $Y_\mu X_\mu$ propagator. After diagonalization of the current eigenstates, one makes the gauge kinetic terms of Eq.(1) diagonal and canonical.

We show in Fig.1 (left) the points that fulfill the WMAP 5σ bound [27] on $\Omega_{DM}$ for $m_{\psi_0} = 10$ GeV in the $(M_{Z_D}, \delta)$ plane. One can clearly see the $Z_D$–pole region when $M_{Z_D} \sim m_{\psi_0}$. One important point is that for a given $M_{Z_D}$ and $m_{\psi_0}$, there exists a unique solution $\delta$ (up to the very small uncertainties at $5\sigma$) fulfilling WMAP constraints : from 3 parameters ($m_{\psi_0}, M_{Z_D}, \delta$), the WMAP constraints reduce it to two ($M_{Z_D}, \delta$).
We show in Fig.1 the points respecting WMAP, and the DAMA/LIBRA (with and without channeling) CoGeNT and CRESST\(^4\) results at 90\% of CL. All the constraints have been calculated for a standard Maxwellian velocity distribution (with mean velocity \(v_0 = 230\) km/s and an escape velocity \(v_{\text{esc}} = 600\) km/s). One can observe in Fig.1 that for all experiments, the regions are quite surprisingly near and correspond to 10 GeV \(M_{\psi} = 30\) GeV and \(10^{-2}\delta10^{-3}\), which is in complete agreement with the measurement of electroweak precision tests. Moreover, such values of \(\delta\) are typical of one loop-order corrections and can easily be generated by heavy-fermions loops in the \(Z - Z_D\) propagator.

3. Anomalies and gamma-ray line

It is well known that any extension of the SM which introduces chiral fermions with respect to gauge fields suffers from anomalies, a phenomenon of breaking of gauge symmetries of the classical theory at one-loop level. Anomalies are responsible for instance for a violation of unitarity and make a theory inconsistent [29, 30]. For this reason if any construction introduces a new fermionic sector to address the DM issue of the SM, it is vital to check the cancelation of anomalies and its consequences on the Lagrangian and couplings. The idea is to add to the Lagrangian local gauge non-invariant terms in the effective action whose gauge variations cancel the anomalous triangle diagrams. There exist two kinds of term which can cancel the mixed \(U(1)_D \times G_{SM}^A\) anomalies, with \(G_{SM}^A\) being one of the SM gauge group \(SU(3) \times SU(2) \times U_Y(1)\): the Chern Simons (CS) term which couples the \(G_{SM}^A\) to the \(U(1)_D\) gauge boson, and the Peccei-Quinn (PQ, or Wess-Zumino (WZ)) term which couples the \(G_{SM}^A\) gauge boson to an axion. In the effective action, these terms are sometimes called Generalized Chern–Simons (GCS) terms [6]:

\[
\mathcal{L}_{\text{inv}} = -\frac{1}{4g_X^2}Y_{\mu
u}Y_{\rho\sigma} - \frac{1}{4g_X^2}X_{\mu\nu}X_{\rho\sigma} - \frac{1}{2}(\partial_{\mu}a_X - M_X X_{\mu})^2 - i\bar{\psi}\gamma^\mu D_{\mu}\psi
\]

\[
\mathcal{L}_{\text{var}} = \frac{C}{24\pi^2}a_X \epsilon^{\mu\nu\rho\sigma}Y_{\mu\nu}Y_{\rho\sigma} + \frac{E}{24\pi^2}a_X \epsilon^{\mu\nu\rho\sigma}X_{\mu\nu}F_{\rho\sigma}. \tag{3}
\]

The Stueckelberg axion \(a_X\) ensures the gauge invariance of the effective Lagrangian and \(g_X\) and \(F_{\mu\nu} = \partial_{\mu}X_{\nu} - \partial_{\nu}X_{\mu}\) are the gauge coupling and field strength of \(U(1)_D\). The axion has a shift transformation under \(U(1)_D\)

\[
\delta X_{\mu} = \partial_{\mu}\alpha, \quad \delta a_X = \alpha M_X. \tag{4}
\]

The \(\mathcal{L}_{\text{var}}\) will generate after the \(SU(2) \times U(1)_Y\) tri-vectorial couplings \(Z_D ZZ\) and \(Z_D Z\gamma\) (\(Z\gamma\) coupling being forbidden by spin- momentum conservation). This generates new annihilation processes \(\psi_0 \psi_0 \rightarrow ZD \rightarrow ZZ/\gamma\) which can be observable through the only monochromatic gamma–ray line with energy \(E_\gamma = m_{\psi_0} \left(1 - \frac{M_Z^2}{4m_{\psi_0}^2} \right)\) [15, 16]. Other models predicts several lines [31, 17], but none of them just one line.

As an illustrative point, we show in the left panel of Fig.2 an example of spectrum from the centre annulus that could be observable by the FERMI telescope, generated by DM annihilation within the pole region respecting WMAP constraint\((m_{\psi_0} = 258\) GeV and \(M_{Z_D} = 591\) GeV\). We can clearly distinguish a \(\gamma\)–ray line centered around \(E_\gamma = m_{\psi_0} \left(1 - \frac{M_Z^2}{4m_{\psi_0}^2} \right)\) above the continuous flux produced by the annihilation process \(\psi_0 \psi_0 \rightarrow ZZ/\gamma\). The expected sensitivy of FERMI telescope after 5 years of data taking is presented in the right panel of Fig.2.

\(^1\) For the CRESST estimation, we used an extrapolation given in the talk of T. Schwetz and the CRESST collaboration [28].

\[\text{VI International Workshop on the Dark side of the Universe (DSU 2010) IOP Publishing}
\text{Journal of Physics: Conference Series 315 (2011) 012014 doi:10.1088/1742-6596/315/1/012014} \]
Figure 2. Left: example of gamma–ray flux respecting WMAP constraint for a DM mass of 258 GeV. Right: monochromatic γ–ray fluxes generated by anomaly-cancelation mechanism in comparison with expected 5σ and 95% CL sensitivity contours (5 years of FERMI operation) for the conventional background and unknown WIMP energy, for an effective scale $\Lambda_X = 1.5$ TeV.

We clearly see in the right panel of Fig.2 that for an effective scale $\Lambda_X = 1.5$ TeV (scale of the “new physics” corresponding to the fermions generating the anomalies), all the parameter space would be observable by FERMI at 95% CL. Indeed, the points that respect the WMAP constraints lie around the pole $M_{ZD} \sim 2m_{\psi_0}$ where $\sim 60\%$ of the annihilation rate is dominated by the $Z\gamma$ final state. This proportion still holds for annihilating DM in the Galactic halo and gives a monochromatic line observable by FERMI.

4. Conclusion
We showed that the existence of a dark $U(1)_D$ gauge sector which interacts with the Standard Model only through its kinetic mixing or anomaly-generated couplings possesses a valid dark matter candidate respecting accelerator, cosmological and the more recent direct detection constraints. Moreover, considering the latest results of DAMA/LIBRA, CoGENT and CRESST, we demonstrated that a specific range of the kinetic mixing ($\delta \sim 10^{-4} - 10^{-3}$) can explain all these excesses for a dark boson mass $M_{ZD} \sim 10–20$ GeV, whereas anomaly cancelation conditions generate a monochromatic γ–ray line from DM annihilation into $Z\gamma$. Such a signature would be a smoking gun signal for these types of constructions It is interesting to notice that other constraints, coming from synchrotron radiation [32] or diffuse gamma-ray emission [33] can give more restriction to the analysis.

Acknowledgements
Y.M. wants to thank particularly E. Dudas, T. Schwetz, G. Belanger, N. Fornengo and A. Romagnoni for useful discussions. The work was supported by the french ANR TAPDMS ANR-09-JCJC-0146 and the spanish MICINNs Consolider-Ingenio 2010 Programme under grant Multi- Dark CSD2009-00064 and the E.C. Research Training Networks under contract MRTN-CT-2006-035505. The author also wants to send all his regards to the "magic" organizing comitee.

References
[1] See e.g. P. Langacker, arXiv:0801.1345 [hep-ph]; M. Pospelov, A. Ritz and M. B. Voloshin, Phys. Lett. B 662 (2008) 53 [arXiv:0711.4866 [hep-ph]]; M. Pospelov, Phys. Rev. D 80 (2009) 095002 [arXiv:0811.1030 [hep-ph]]; M. Pospelov, A. Ritz and M. B. Voloshin, Phys. Lett. B 662 (2008) 53 [arXiv:0711.4866 [hep-ph]]; J. Kumar and J. D. Wells, Phys. Rev. D 74, 115017 (2006) [arXiv:hep-ph/0606183]; B. Batell,
ph/0703512].

[32] C. Boehm, J. Silk and T. Ensslin, arXiv:1008.5175 [astro-ph.GA]; D. Hooper and T. Linden, arXiv:1011.4520 [astro-ph.HE].

[33] C. Arina and M. H. G. Tytgat, arXiv:1007.2765 [astro-ph.CO].