Pseudoscalar mesons with symmetric bound state vertex functions on the light front

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We study the electromagnetic form factors, decay constants and charge radii of the pion and kaon within the framework of light-front field theory formalism where we use an ansatz for the quark-meson interaction bound-state function which is symmetric under exchange of quark and antiquark momentum. The above mentioned observables are evaluated for the + component of the electromagnetic current, $J^+$, in the Breit frame. We also check the invariance of these observables in other frames, whereby both the valance and the non-valence contributions have to be taken into account, and study the sensitivity of the electromagnetic form factors and charge radius to the model’s parameters; namely, the quark masses, $m_u = m_d$, $m_s$, and the regulator mass, $m_R$. It is found that after a fine tuning of the regulator mass, i.e. $m_R = 0.6$ GeV, the model is suitable to fit the available experimental data within the theoretical uncertainties of both the pion and kaon.

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I. INTRODUCTION

The theory of strong interactions, Quantum Chromodynamics (QCD), has been the object of theoretical and experimental scrutiny for four decades now and is perturbatively well defined. Indeed, at large momentum transfer, perturbative calculations successfully describe a wealth of subatomic phenomena. However, QCD is also a theory whose elementary excitations are confined and nonperturbatively QCD may well be rigorously defined but a full solution is yet out of reach: there exists no simple Schrödinger picture of a many-body Hamiltonian model’s parameters; namely, the quark masses, $m_u = m_d$, $m_s$, and the regulator mass, $m_R$. It is found that after a fine tuning of the regulator mass, i.e. $m_R = 0.6$ GeV, the model is suitable to fit the available experimental data within the theoretical uncertainties of both the pion and kaon.

Thus, nonperturbatively QCD may well be rigorously defined but a full solution is yet out of reach: there exists no simple Schrödinger picture of a many-body Hamiltonian model’s parameters; namely, the quark masses, $m_u = m_d$, $m_s$, and the regulator mass, $m_R$. It is found that after a fine tuning of the regulator mass, i.e. $m_R = 0.6$ GeV, the model is suitable to fit the available experimental data within the theoretical uncertainties of both the pion and kaon.

On the other hand, in understanding the dynamical properties of nonperturbative QCD, the light pseudoscalar mesons and in particular the pion play a crucial role. Remarkably, the latter is a bound state of massive antiquark-quark pairs as well as the almost massless Goldstone boson associated with chiral symmetry breaking. There have been many studies of their static properties [7–12] and their dynamical properties have also been investigated theoretically [13–17] and experimentally [64–74].

Taking advantage of the simple structure of the Fock space and the vacuum in light-front quantization, various hadronic properties of bound states, such as decay constants and electromagnetic form factors of the pion, kaon and nucleon, have been calculated [16, 17, 52–62] and successfully compared with their experimental values [64–74]. Since the light front component, $J^+$, has been successfully employed to calculate electromagnetic form factors [7, 31, 63, 75–79], the light-front approach also offers a theoretical framework to extract from them useful information on the valence and non-valence components of the meson’s wave function.

In the present simultaneous study of electromagnetic form factors, charge radii and decay constants, we adopt the light-front field theory formalism of Refs. [7, 15] wherein the Bethe-Salpeter amplitude of the $qar{q}$ bound states was modeled for two different momentum constellations, namely a symmetric [17] and nonsymmetric vertex model [16]. Here, the vertex refers to the $qar{q}$ pair coupling to the pseudoscalar meson in an effective Lagrangian. Using a nonsymmetric vertex model, E. O. Silva et. al. [58] calculated the aforementioned

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pion and kaon observables which are in agreement with experimental data. However, a momentum distribution of the meson that is symmetric under the exchange of the quark and antiquark momenta is more realistic and such a model for the Bethe-Salpeter amplitude should improve the description of or at least equally well reproduce all observables presented in Ref. [58]. Thus, we here use the same component, \( J^+ \), of the light-front electromagnetic current, though with a symmetric momentum description of quark-meson bound-state vertex. It is important to recall here that the choice of \( J^+ \) with the Drell-Yan condition \( q^+ = 0 \) guarantees that pair-term contributions (non-valence terms) vanish \([16, 17]\). On the other hand, to preserve rotational symmetry, the pair contribution must be included \([18–23]\). Consequently, we here employ both the valence and non-valence contributions considering the case \( q^+ \neq 0 \).

This paper is organized as follows: Sec. II serves to summarize the general framework where subsequently the different physical observables, namely the electromagnetic form factors, decay constant and charge radii are discussed in turn. In Sec. III, we present our numerical results and analyze the observables' dependence on variation of the model parameters. In the last section, we give our conclusions.

II. THE MODEL

In this section, we briefly summarize the model and the computational tools of the light-front formalism required to investigate the pseudoscalar meson's electromagnetic form factors, charge radii and the decay constants. Our approach is based on similar earlier work \([16, 17]\), where the following effective Lagrangian for the \( \bar{q}q \) bound state was employed:

\[
E_{\text{eff}} = -ig \bar{\phi} \gamma^5 \phi \bar{q} q ,
\]

where \( g = m_0 - /f_0 - \) is the coupling constant, \( m_0 - \) and \( f_0 \) denote the mass and decay constant of a pseudoscalar meson, respectively, and \( \bar{\phi} \) represents the scalar field. We make a symmetric ansatz for the \( \bar{q}q \)-meson vertex which describes the bound state,

\[
\Lambda(k, P) = C [(k^2 - m_R^2 + i\epsilon)^{-1} + ((P - k)^2 - m_R^2 + i\epsilon)^{-1}],
\]

where it is clear that the \( \Lambda(k, P) \) is symmetric under the exchange of the quark and antiquark momenta, \( k \) and \( (P - k) \); \( P \) is the total momentum of the meson and \( C \) a normalization constant. In the following, we discuss the analytic light-front formulation of the electromagnetic form factor, charge radius and decay constant.

A. Electromagnetic form factors

The covariant electromagnetic form factor, \( F_{\mu}^{\text{em}} \), is defined by a matrix element where the electromagnetic current, \( J_\mu = e_q \bar{q} \gamma_\mu q \), is sandwiched between the initial and final bound states of the same meson:

\[
P_\mu F_{\mu}^{\text{em}}(q^2) = \langle M_0 - (p')|J_\mu|M_0 - (p)\rangle ,
\]

where \( M_0 - = \pi^+, K^+ \), \( P_\mu = (p + p')_\mu \) and \( q^2 = (p' - p)^2 \) is the square of the momentum transfer.

The electromagnetic form factor in the impulse approximation is obtained from triangle diagrams, each of which contains one spectator quark. In this approximation, the covariant electromagnetic current of a pseudoscalar meson, \( J_\mu \), that enters Eq. (3) can be written as follows \([80, 81]\):

\[
J_\mu = N \int \frac{d^3k}{(2\pi)^3} \text{Tr} \left[ \frac{1}{k - m_\bar{q} + i\epsilon} \gamma^5 \right] \frac{1}{k - p' - m_q + i\epsilon} \gamma_\mu \Lambda(k, p') \Lambda(k, p) + [q \leftrightarrow \bar{q}] ,
\]

with the normalization,

\[
N = -\frac{2i e_q \hat{m}_0^2 - N_c}{f_0'}, \quad \hat{m}_0 - : \frac{m_q + m_\bar{q}}{2} ,
\]

where \( N_c = 3 \) is the number of colors and \( f_0 \) the pseudoscalar weak decay constant.

The light front-form variables are,

\[
k^+ = k^0 + k^3 \quad k^- = k^0 - k^3 \quad \vec{k}_\perp \equiv (k^1, k^2)
\]

\[
q^+ = \sqrt{-q^2} \sin \alpha, \quad q_\perp = \sqrt{-q^2} \cos \alpha, \quad q_\perp = 0
\]

\[
q^2 = q^+ q^- - \vec{q}_\perp^2 .
\]

Here, we use the Drell-Yan condition, \( q^+ = 0 \), in the Breit frame, which implies \( \alpha = 0 \). However, our results are frame invariant, i.e. invariant for \( \alpha \neq 0 \) where both the valence and the non-valence contributions become important. After introducing the front-form variables in Eq. (5) and using \( \gamma^+ = \gamma^0 + \gamma^3 \) to obtain the \( J^+ \) component of the current in Eq. (4), the electromagnetic form factor becomes:

\[
F_{\mu}^{\text{em}}(q^2) = \frac{N}{P^\perp} \int \frac{d^2k_\perp}{(2\pi)^2} dk^+ dk^- \text{Tr} [\Gamma(O^+)] \times \Gamma(k^+, p^+, p'^+) + [q \leftrightarrow \bar{q}] ,
\]

In Eq. (6), the trace in light-front coordinates is,

\[
\frac{1}{4} \text{Tr}[\Gamma(O^+)] = \frac{1}{4} k^+ q^+ + (k^+ - p^+ - p'^+) (k^+ - k^-)
\]

\[
- k^- p^+ p'^+ - (p^+ k_\perp \cdot p_\perp + p^+ k_\perp \cdot p'_\perp)
\]

\[
- k^+ (2k_\perp^2 + m_\bar{q}^2 - 2m_q m_\bar{q}) - (p^+ + p'^+) m_q m_\bar{q} ,
\]

and

\[
\Gamma(k^+, p^+, p'^+) = \frac{\Lambda(k^+, p^+) \Lambda(k^+, p'^+)}{(k^2 - m_\bar{q}^2 + i\epsilon)(p^2 - k^2 - m_q^2 + i\epsilon)} \times \frac{1}{((p' - k)^2 - m_\bar{q}^2 + i\epsilon)} ,
\]
where \( k^2 = k^+(k^- - k_{on}^-) \) where \( k_{on}^- \) is the on-energy-shell value of the corresponding momentum given by,

\[
k_{on}^- = \frac{k_+^2 + m_0^2}{k^+}.
\] (9)

In terms of light-front variables, the bound-state function in Eq. (2) becomes,

\[
\Lambda(k^+, p^+) = \frac{C}{(p^+ - k^+)(p^- - k^- - (p-k)^2 + m_0^2 - i\epsilon/p^+ - k^+)} + \frac{C}{k^-(k^- - k_{on}^- + m_0^2 - i\epsilon/k^+)}.
\] (10)

Collecting all ingredients from Eqs. (7)-(10), we insert them in Eq. (6) and after \( k^- \) energy integration (see appendix of Ref. [17]) with \( x = \frac{k^-}{p^-} \) the electromagnetic form factor can be rewritten,

\[
F_{0-}^{em} = \frac{N}{p^+} \int \frac{d^2 k_\perp dx}{x(1-x)} \Phi^*(x, k_\perp) \Phi(x, k_\perp) \theta(x) \theta(1-x) \times \left[ \frac{1}{2} k^+ q_\perp^2 - k_{on}^- p^+ p^+ - (p^+ k_\perp \cdot p_\perp + p^+ k_\perp \cdot p'_\perp) \right]
\] (11)

where \( N = \frac{N_c^2}{(2\pi)^3} \) and

\[
\Phi(x, k_\perp, p^+, p'_\perp) = \left[ \frac{1}{(1-x)(m_0^2 - M^2(m_0^2, m_R^2))} \right] + \frac{1}{x(m_0^2 - M^2(m_R^2, m_0^2))} \left[ 1 \right] + \frac{1}{(1-x)(m_0^2 - M^2(m_0^2, m_0^2))}, + [q \leftrightarrow \bar{q}]
\] (12)

with

\[
M^2(m_a^2, m_b^2) = \frac{k_+^2}{x} + \frac{(p-k)^2}{1-x} + m_0^2.
\]

Note that the appearance of the second term of Eq. (12) is due to the symmetric character of the meson-quarks vertex, which is absent in Refs. [16, 58] where the authors consider a nonsymmetric behavior of the vertex function. As it was shown [16, 17, 25], to preserve general covariance the non-valence contribution is mandatory. Thus, in Eq. (11) the step functions \( \theta(x) \) and \( \theta(1-x) \) delimit the integration interval, \( 0 < k < p^+ \), of the valence contribution, whereas the interval, \( p^+ < k^- < p^+ \), corresponds to the non-valence contributions to electromagnetic current [17].

### B. Charge radius and decay constant

The mean-square electric charge radius of a meson is a relevant quantity and correlated with the electromagnetic form factor,

\[
F_{0-}^{em}(q^2) \simeq 1 - \frac{1}{6}(r_{0-}^2)q^2.
\] (13)

Differentiation with respect to \( q^2 \), of the above equation yields the charge radius,

\[
\langle r_{0-}^2 \rangle = \frac{dF_{0-}^{em}}{dq^2} \bigg|_{q^2=0}.
\] (14)

A relevant observable and also our main constraint on the model’s parameters is given by the weak decay constant, \( f_{0-} \). The decay constant of a \( q\bar{q} \) bound state can be found from the following matrix element of the partially conserved axial-vector current:

\[
\langle 0|A^\mu|0^- \rangle = i f_{0-} \gamma^\mu q.
\] (15)

where \( A^\mu = \bar{q} \gamma^\mu \gamma^5 q \), is the axialvector current. The weak decay constant is given by,

\[
f_{0-} = \frac{ieN_c}{f_{0-}} \int \frac{d^4k}{(2\pi)^3} \text{Tr} \left[ \frac{\gamma^5}{k^2 - m_0^2 + i\epsilon} + \gamma^5 \frac{(k - p)}{(k - p)^2 - m_0^2 + i\epsilon} \right] \Lambda(k, p)
\] (16)

We make use of the + component of the axialvector current \( A^\mu \) and after integration over \( k^- \), one obtains the decay constant in terms of the valence component of the model:

\[
f_{0-} = \langle 0|A^\mu|0^- \rangle f_{0-} = \frac{\sqrt{N_c}}{4\pi^3} \int \frac{d^2 k_\perp dx}{x} (4m_0 + 4m_q(1-x)) \Phi(x, k_\perp, m_{\pi}, \bar{q}).
\] (17)

### III. NUMERICAL RESULTS AND DISCUSSION

The model introduced in Sec. II contains three free parameters, namely, the regulator mass, \( m_R \), and the two constituent quark masses, \( m_a = m_d \) and \( m_s \), where the strange quark mass is taken from the study in Ref. [24]. The main focus of this study is to constrain the parameters of a more realistic bound-state ansatz to accommodate the available experimental data on the pion and kaon elastic form factors, decay constants and charge radii. In addition, it is also instructive and important to check the explicit dependence of these observables on the model’s parameter.

Regarding these goals, we know from a previous study [17] that the value of the regulator mass, \( m_R = 0.6 \) GeV, reproduces well all experimental data on the pion observables mentioned above. It is worthwhile to check whether \( m_R = 0.6 \) GeV is also consistent with the kaon data, for which we compute the values of the decay constants and charge radii and compare them with their experimental values.

The calculated values of the observables listed in Table 1 show that \( m_R = 0.6 \) GeV is also a suitable value for the kaon. Moreover, for \( m_s = 0.44 \) GeV,
seen, the elastic form factors in both figures are monotonically decreasing functions of $q^2$ with increasing hardness for larger values of $m_R$. Fig. 1 also informs us that the available experimental pion data lie in the interval of $0.1 \text{ GeV} \leq m_R \leq 1 \text{ GeV}$ and from Fig. 2 we deduce that the kaon’s experimental form factor data are better reproduced for $m_R \gtrsim 0.5 \text{ GeV}$ which coincides with our privileged value, $m_q = 0.6 \text{ GeV}$.

Similarly, in Figs. 3 and 4, the electromagnetic form factors of the pion and kaon are plotted as a function of $q^2$ for different values of $m_u = m_d$ whereas $m_R = 0.6 \text{ GeV}$ is fixed. One observes a likewise behavior of the electromagnetic form factors, i.e. $F_{\text{em}}(q^2)$ becomes harder for increasing values of $m_u = m_d$. We note that for the symmetric vertex function with $m_R = 0.6 \text{ GeV}$, the light-quark mass should be in the range $0.15 \lesssim m_q \lesssim 0.5 \text{ GeV}$, where the constituent quark mass, $m_q = 0.22 \text{ GeV}$, appears to be the most favorable value to accommodate the experimental data.

Next, the explicit dependence of the charge radii $\langle r_\pi \rangle$

\begin{table}[h]
\centering
\begin{tabular}{|c|c|c|c|c|}
\hline
Decay Constant and Charge radius & $m_R = 0.6 \text{ GeV}; m_s = 0.44 \text{ GeV}$ & $m_R = 0.6 \text{ GeV}; m_s = 0.44 \text{ GeV}$ \\
\hline
Pion & $m_u = m_d = 0.22 \text{ GeV}$ & $m_u = m_d = 0.25 \text{ GeV}$ & $m_u = m_d = 0.22 \text{ GeV}$ & $m_u = m_d = 0.25 \text{ GeV}$ \\
\hline
$f_0^-$ & 93.12 MeV & 101.85 MeV & 110.81 MeV & 113.74 MeV \\
$\langle r_0^- \rangle$ & 0.736 fm & 0.670 fm & 0.754 fm & 0.687 fm \\
\hline
\end{tabular}
\caption{Calculated decay constants and charge radius for two light-quark masses and corresponding experimental values. The experimental date with errors bar, are; $f_\pi^{\exp} = 92.42 \pm 0.021 \text{ MeV}$, $\langle r_\pi \rangle^{\exp} = 0.672 \pm 0.008 \text{ fm}$, $f_K^{\exp} = 110.38 \pm 0.1413$ and $\langle r_K \rangle^{\exp} = 0.560 \pm 0.031 \text{ fm}$. Experimental date from [64, 70, 82].}
\end{table}

\begin{figure}[h]
\centering
\includegraphics[width=0.45\textwidth]{fig1}
\caption{The electromagnetic form factor of the pion as a function of space-like $q^2$. The curves correspond to the different values of $m_R$ with fixed quark masses: $m_u = m_d = 0.220 \text{ GeV}$. Experimental data: Ref. [66] (circle), Ref. [65] (square), Ref. [68] (diamonds), Ref. [69] (up triangle), Ref. [73] (down triangle).}
\end{figure}

\begin{figure}[h]
\centering
\includegraphics[width=0.45\textwidth]{fig2}
\caption{The electromagnetic form factor of the kaon as a function of space-like $q^2$. The curves correspond to the different values of $m_R$ with fixed quark masses: $m_u = m_d = 0.220 \text{ GeV}, m_s = 0.44 \text{ GeV}$. Experimental data: Ref. [83] (square), Ref. [64] (circle).}
\end{figure}
FIG. 3: The electromagnetic form factor of the pion as a function of space-like $q^2$. The curves correspond to different values of $m_q = m_u$ with regulator mass $m_R = 0.60$ GeV. Experimental data: Ref. [66] (circle), Ref. [65] (square), Ref. [68] (diamonds), Ref. [69] (up triangle), Ref. [73] (down triangle).

and $(r_K)$ on $m_u = m_d$ is depicted in Fig. 5 from which we deduce that the charge radius exhibits a nonlinear behavior and decreases strongly for large values of $m_u = m_d$, as expected. For instance, for $m_u = m_d = 0.15$ GeV, the size of the pion (kaon) charge radius is about 0.98 fm ($\sim$ 1 fm), whereas for $m_u = m_d = 0.5$ GeV this size reduces to 0.43 fm (0.45 fm). A similar behavior can be seen in Fig. 6, where the charge radii are plotted against the regulator mass $m_R$ with $m_u = m_d = 0.22$ GeV and $m_z = 0.44$ GeV fixed. The solid circle [70] and the square [64] are the experimental charge radii values of the pion and kaon, respectively.

FIG. 4: The electromagnetic form factor of the kaon as a function of space-like $q^2$. The curves correspond to different values of $m_q = m_u$ with fixed masses: $m_s = 0.44$ GeV, $m_R = 0.60$ GeV. Experimental data: solid circle [64] (circle) and square [63].

FIG. 5: Charge radii $(r_{0-})$ of the pion and kaon as a function of the constituent quark mass $m_u = m_d$ with $m_s = 0.44$ GeV and fixed regulator mass $m_R = 0.6$ GeV. The solid circle [70] and square [64] are the experimental values for the charge radii of the pion and kaon, respectively.

and $f_K$, are plotted in Fig. 7 and Fig. 8 as functions of $m_u = m_d$.
FIG. 7: The weak decay constants \( f_{0^-} \) of the pion (solid curve with \( m_u = m_d = 0.22 \) GeV) and kaon (dashed curve: \( m_s = 0.44 \) GeV) as a function of \( m_u = m_d \); dotted curve: kaon decay constant as a function of \( m_s \). In all curves \( m_R = 0.6 \) GeV. The solid circle and square are the experimental decay constants of the pion and kaon, respectively [82].

\((m_R = 0.6 \) GeV) and \( m_R (m_u = m_d = 0.22 \) GeV), where in both cases \( m_s = 0.44 \) GeV. We observe that in contrast to the charge radii the decay constants are continuous increasing functions of \( m_q \) and \( m_R \). Moreover, the charge radii and the decay constants are more sensitive to the quark mass values, \( m_q \), than to \( m_R \). It is worthwhile to point out, as discussed in Ref. [58] for the non-symmetric vertex, that the sensitivity of the kaon decay constant to the strange quark mass is very modest, whereas for the present symmetric vertex the \( m_s \) dependence is quite significant.

We stress that the regulator mass \( m_R = 0.6 \) yields the best fit to the experimental values of the decay constants and charge radii as discussed above. Furthermore, the decreasing (increasing) behavior of the charge radii (decay constants) with \( m_q \), as depicted in Figs. 5 and 7, satisfies Tarrach’s relation [84], i.e. \( \langle r_{0^-} \rangle \sim 1/m_q \) and \( f_{0^-} \sim 1/\langle r_{0^-} \rangle \).

IV. CONCLUSIONS

We reexamined the light-front approach to the light pseudoscalar mesons [16, 17, 58] by considering a symmetric \( \bar{q}q \) bound-state function [17]. In this framework and with the given symmetric ansatz for the bound-state vertex function, we calculated the charge radii, \( \langle r_{\pi} \rangle \), and \( \langle r_K \rangle \), weak decay constants, \( f_{\pi} \) and \( f_K \) and the electric form factors, \( F_{\pi}^{em}(q^2) \) and \( F_{K}^{em}(q^2) \).

To constrain our model parameters, namely, the quark masses \( m_u, m_d \) and \( m_s \) and the regulator mass \( m_R \), we first adjusted their values to reproduce the experimental weak decay constants. In doing so, we imposed the same regulator mass for both the pion and kaon and found that \( m_R = 0.6 \) GeV is a suitable value to describe all experimental data on \( F_{\pi(K)}^{em}, \langle r_{\pi(K)} \rangle \) and \( F_{\pi(K)} \) within the reasonable theoretical uncertainties. The numerical results also show this model significantly breaks down for \( m_R \geq 1 \) GeV, which was already demonstrated in the the case of a nonsymmetric vertex function [58].

Moreover, the explicit dependence of charge radii and decay constants on the quark masses, \( m_u = m_d \) and \( m_s \), not only satisfies Tarrach’s relation but also favors the range of mass values commonly chosen within the light-front model. In addition, by using the privileged values of the model’s parameters \( m_u = m_d = 0.22 \) GeV, \( m_s = 0.44 \) GeV and \( m_R = 0.6 \) GeV, we find that the pion-to-kaon decay constant ratio is in excellent agreement with its experimental value, i.e. the mismatch is barely 0.67%. Lastly, the present numerical investigation suggests that these parameter values could also be used to study dynamical properties of other pseudoscalar and vector mesons or may apply to heavy-to-light transition form factors.
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