Inverse spin switch effects in Ferromagnet / Superconductor hybrids with strong ferromagnets

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(Dated: March 23, 2022)

In F/S/F trilayers where the magnetization directions of the F layers can be controlled separately, it has theoretically been predicted that the antiparallel (AP) configuration can have a higher superconducting transition temperature \( T_c \) than the parallel configuration. This is the so-called spin switch, which also experimentally has been found for the case of weak ferromagnets. Here we show that strong ferromagnets yield the opposite effect. We study the transport properties of F/S/F trilayers with \( F = \text{Ni}_{0.80}\text{Fe}_{0.20} \) (Permalloy, Py) and \( S = \text{Nb}, \) structured in strips of different sizes. Using two different thicknesses for the Py layers, we can switch in a well-defined way between the AP- and P-configurations. In the superconducting transition we find a clear increase of the resistance in the AP-state. We ascribe this to enhanced reflection of spin-polarized quasiparticles at the S/F interfaces which leads to a stronger suppression of superconductivity on the S-side.

PACS numbers: 74.45.+c, 74.78.-w, 85.25.-j

One of the interesting phenomena which is currently searched for in hybrids of superconductors (S) and ferromagnets (F) is the so-called superconducting spin-switch effect. Basically, the effect can occur in F/S/F trilayers in which the direction of the magnetization of one F-layer can be varied with respect to the other. It was predicted some time ago that, in a current-in-plane (CIP) geometry, and for a thickness of the S-layer \( d_S \) of the order of its superconducting coherence length \( \xi_S, \) the transition temperature \( T^{AP}_c \) in the antiparallel (AP) state is higher than the one in the parallel (P) state, \( T^{P}_c. \) For particular choices of the different layer thicknesses it should even be possible to find full re-entrant behavior, controlled with only a small switching field \( h_{Pc}. \) The effect is reminiscent of F/N/F spin valves (N being a normal metal), with one important difference. In the normal metal spin valves, the resistance is lowest in the parallel configuration, since in terms of a two-spin-current model, it is determined by the spin channel with the smallest resistance. In the superconducting case, the antiparallel configuration yields the lowest (zero) resistance since the Cooper pair samples opposite exchange fields, which are less pair-breaking than parallel fields.

Full re-entrant behavior has not yet been observed in superconducting spin valves. Two experiments were reported. Both were with a device consisting of weakly ferromagnetic CuNi and superconducting Nb, and in both cases the reported \( T^{AP}_c \) was only about 5 mK or 2.5 mK \( h_{Pc}. \) Here we want to show that high spin polarization actually leads to the opposite effect. In trilayer combinations of ferromagnetic Permalloy (Ni_{0.80}Fe_{0.20}; Py) with superconducting Nb, we find that, in the superconducting transition, switching from the P to the AP configuration leads to an increase rather than a decrease of the resistance when measured in a CIP geometry. The effect is strong: it is very similar to recently reported findings on trilayers of ferromagnetic La_{0.7}Ca_{0.3}MnO_{3} (L), where \( P_L \) is expected to be close to 100 \% and superconducting YBa_{2}Cu_{3}O_{7} (Y), and we offer a similar explanation in terms of reflection of spin-polarized quasiparticles near the interface. Where the conclusions in ref. \( c \) had to be based on indirect evidence for the AP configuration, the use of Py allows to demonstrate switching effects directly. Together, the experiments based on oxides and on simple metals show that the effect is generic for the limit of high spin polarization.

Samples of Nb/Py were prepared by sputter deposition in an ultrahigh vacuum system. Thick Nb films have a \( T_c \) of 9.2 K, similar to the bulk. From the upper critical fields we extract a value for the Ginzburg-Landau coherence length \( \xi_{GL}(0) \approx 13 \text{ nm}. \) Using \( \xi_{GL}(0) = 0.86\sqrt{\xi_{GL}(0)} \), with the BCS coherence length \( \xi_0 \approx 40 \text{ nm}, \) this yields a value for the normal state elastic mean free path \( \ell_N \approx 5.5 \text{ nm}. \) Samples were structured by e-beam lithography in bridges of 0.5 mm \( \times \) 4 mm (‘large’ samples) or in bridges of 3 \( \mu \text{m} \times 20 \mu \text{m} \) (‘small’ samples). For both large and small samples we used a design in which the contacts were included in the geometry and simple bars with gold contacts in order to minimize problems with stray fields from contact pads or arms. Ferromagnetic Py possesses a large spin polarization (45 \%), but also shows well-defined magnetization switching at low fields. Care was taken to align the long axis of the bars with the easy axis of magnetization \( \hat{e}_z, \) which is induced by the residual magnetic fields in the sputtering machine. Magnetic fields were applied in the plane of...
the sample, along the bars and therefore along \( \hat{e}_x \). We also made use of the fact that the coercive field \( H_c \) of the Py layers depends on their morphology as well as on thickness. We consistently find that a thicker Py layer deposited on the substrate has a lower value of \( H_c \) than a thinner layer deposited on top of the Nb layer. In this way it is quite possible to have well-defined P- and AP-regimes. Typically, we used 50 nm for the Py bottom layer and 20 nm for the Py top layer, which yields coercive or switching fields around 1 mT - 2 mT (50 nm) and 8 mT - 10 mT (20 nm). For the Nb layer the smallest thickness \( d_{Nb} \) was around 25 nm, which yields transition temperatures around 4 K and is already close to the critical thickness \( d_c \) for the trilayer.

Fig. 1 shows a compilation of different measurements on a sample \( s/Py(50)/Nb(25)/Py(20) \) (with \( s \) denoting the substrate and the numbers the layer thickness in nm). The magnetization was measured at 5 K on the unstructured sample by SQUID magnetometry. The switching of the layers is well defined and from the magnitude of the jumps it can be seen that the 50 nm layer switches at \( \pm H_{c,50} \approx 1.5 \) mT, and the 20 nm layer at \( \pm H_{c,20} \approx 9.5 \) mT, leaving a large field range for the AP-state. For transport measurements, the sample was structured as a large bar with gold contacts and showed a resistive transition around 3.7 K with a width \( \Delta_T \) of 100 mK. The resistance \( R \) was measured as function of the in-plane field \( H_a \) at a temperature of 3.66 K, as shown in Fig. 1. Starting at high fields, \( R \) decreases until \( +H_{c,20} \) where it starts to rise slowly. At \( -H_{c,50} \) a small but clear upward jump occurs. This is the field where the alignment of the Py layers becomes AP. In this regime \( R \) rises further to a peak, followed by a steep decrease to a dip at around \( -H_{c,20} \). Now the sample is in the P state and \( R \) starts to rise slowly again. The behavior is mirrored in increasing fields. The strong peaks in the resistance therefore appear to be connected to the AP alignment, just as in the case of the L/Y/L trilayers of ref. 8. The dips at \( \pm H_{c,20} \) are well known and are produced by the magnetic domains which occur around the Py coercive field in the 20 nm layer. They can also be found in Py/Nb bilayers, and are due to a lower averaged exchange field sensed by the Cooper pair, as we have demonstrated previously 4. For the 50 nm layer they also should be present, and a small dip is actually observed in backward sweep at \( +1.5 \) mT, but it is masked by the increase of resistance resulting from the P \( \rightarrow \) AP switch. The effect of the domain state on \( T_c \) is shown in the inset, where \( R(T) \) is given at two different fields: at 5 mT coming from high field (sample in the P-state), and at 10 mT coming from low fields (sample in the domain state). The difference in the temperature where zero resistance is reached is at most 30 mK, similar to the earlier findings. Finally, the rise in \( R \) at \( +H_{c,20} \) when coming from high fields may appear puzzling, since no switching takes place at this field. Domain formation, however, already does set in: close inspection of M(H) shows that decrease already starts before \( H_{c,50} \) is reached, and this should lead to small amounts of AP orientations.

The data therefore suggest that the AP state shows larger resistance than the P state, but for the large sample the behavior is sluggish because of domain effects. Next we consider some much smaller samples. Fig. 2 shows data on a sample \( s/Py(50)/Nb(26)/Py(20) \) with a bridge of 3 \( \mu m \times 20 \) \( \mu m \) and contacts included. The transition (see inset) is quite broad, \( \Delta_T \approx 600 \) mK and the measurement is taken at 3.80 K, close to the onset at 4 K. After correcting for a small offset field, the values for \( \pm H_{c,50} \) and \( \pm H_{c,20} \) are 2.7 mT and 10 mT, respectively. The switching behavior is now perfectly well defined. Also noticeable is the absence of the dips which we ascribed to the domains. This is again in agreement with our earlier observations and due to the fact that no stable domain state is formed in these
small samples during the switching \cite{Gu2019}. Samples with similar thicknesses show the same behavior, although the switching is not always a perfect one-step process; sometimes, several steps (both up and down) can be seen. This may not be surprising, since the properties of Py, and the direction of the easy axis are very sensitive to the preparation conditions. Next we increase $d_{Nb}$. The transition width now gradually decreases. Fig. 4 shows data on a sample s/Py(50)/Nb(60)/Py(20) with $\Delta_T \approx 100$ mK and a normal state resistance of 9.89 $\Omega$. In this case, the sample was a simple bar of 3 $\mu$m x 20 $\mu$m with Au contacts, which we show to make clear that the effect can be found with different contact geometries. The behavior of $R(H_a)$ is shown at $T = 7.46$ K, halfway the transition, and at $T = 7.40$ K, close to the bottom. The switching behavior is still sharp and clear, quite similar to the previous sample. Note that the size of the resistance variation has become much smaller at the lower temperature. For samples with $d_{Nb} = 80, 90$ nm and $\Delta_T \leq 50$ mK, the effect became very small. Above $T_c, R(H_a)$ of the samples only shows tiny dips around $\pm H_{c2,50}$, which are due to the anisotropic magnetoresistance (AMR) effect for configurations where the current is parallel to the applied magnetic field \cite{Mourigal2016}.

The observed effects are therefore due to the onset of the superconducting state. For all geometries, in the resistive transition, the AP state has a higher resistance than the P state, which is opposite to the effects predicted \cite{Blonder1982, Tinkham2004} and observed for weak ferromagnets \cite{Klapwijk1985, Gu2019}, but similar to the observations of ref. \cite{Yamashita2017}. In terms of the magnetoresistance $MR = \Delta R = (R_{AP} - R_P)/R_P$, the effect becomes quite large due to the decreasing value of $R$, but for the physics it is more relevant to note that in terms of the normal-state resistance of the samples, it is a fraction of the order of 5% - 10%, at least in the upper half of the transition. It is then of interest to compare the results of these CIP measurements to the results of current-perpendicular to-plane (CPP) measurements on stacks of Py/Nb/Py, performed by Gu et al. \cite{Gu2018}, for different thicknesses $d_{Nb}$ of the Nb layer in a range between 30 nm and 100 nm. In the CPP-case, a small (1%) positive ($R_{AP} > R_P$) effect was present in the normal state, which persisted below $T_c$. However, it was found to decrease with decreasing temperature, which is different from our CIP data which show an initial increase of MR (from 0 in the normal state). A basic explanation of the CPP data was given in terms of the diffusion of spin-polarized quasiparticles (qp) with energies $E_{qp}$ below the gap. The MR effect in the normal state is due to the standard mechanism of increased spin scattering in the AP configuration. This becomes smaller through spin memory loss controlled by the spin diffusion length $\ell_{sd}$, but when the intermediate layer is a superconductor, where spin is carried only by quasiparticles and not by Cooper pairs, MR will also decrease due to the loss of quasiparticles to the condensate, which is controlled by the qp diffusion length $\xi_{ap}$ \cite{Glazman2017}. From the temperature dependence of $\Delta R$, a value of $\xi_{ap} \approx 16.5$ nm was found, very close to $\xi_{GL}(0) \approx 13$ nm of the Nb. This is actually not surprising. It was already shown by Blonder, Tinkham and Klapwijk that the characteristic decay of evanescent quasiparticles with energies inside the gap (the current-to-supercurrent conversion length) is given by $1.22 \xi_{GL}(T)$ \cite{Blonder1982}. The number also shows that spin loss by spin scattering only plays a minor role, since $\xi_{ap} \approx 50$ nm \cite{Glazman2017}. This analysis was confirmed in recent theoretical work by Yamashita et al. \cite{Yamashita2017}, who considered Andreev reflections and direct transmission of spin-polarized quasiparticles in F/S/F systems. They found slightly larger values for $\xi_{ap}$, and also that $\xi_{ap}$ decreases with increasing $d_{Nb}$. This was ascribed to the proximity effect in this all-metal system, which suppresses the average gap as long as $d_{Nb}$ is comparable to $\xi_{GL}$.

Translating this description to our CIP case, the first thing to note is the differences in MR in the normal state. This is easily explained. In CPP, for diffusive systems, the dependence of $\Delta R/R$ on the spacer thickness $d_N$ is $\propto e^{-d_N/\ell_{sd}}$ \cite{Glazman2017}. If MR effects are present in the Py/Nb/Py system, they can be witnessed for the range of thicknesses used in the CPP experiment. On the other hand, in CIP the attenuation is $\propto e^{-d_N/\ell_N}$, with $\ell_N$ the elastic mean free path of the normal metal \cite{Glazman2017, Glazman2017}. For our Nb, we estimated before that $\ell_N \approx 5.5$ nm, much smaller than the spacer thickness of 25 nm. So, if MR effects are found in CPP, they will not be observed in CIP in the same thickness range. This changes in the superconducting transition, since then quasiparticles appear with a much longer range because of the divergence close to $T_c$. We now offer a similar line of reasoning as ref. \cite{Glazman2017}. Current is flowing in the plane of the films, there is no voltage difference perpendicular to the layers, so electrons scattering out of the F layer turn into low-energy spin-polarized quasiparticles, which can diffuse across the S layer. In the AP configuration, a larger number experiences reflection at the other interface than in the P configuration. Although the spin determines this reflection process, the result should probably not be called spin accumulation since there is...
no net charge or spin transport through the interface in contrast to the CPP case. The reflection however also leads to a larger number of quasiparticles on the S-side for the AP case than for the P case, and this translates (selfconsistently) into a gap suppression on the S-side. This gap suppression is observed as an increased resistance. Based on the magnitude of the switching, we can estimate the size of the layer where this takes place for temperatures close to \( T_c \). The reflection however leads to a larger number of quasiparticles on the S-side for the AP case than for the P case, and this also leads to a larger number of quasiparticles on the S-side. This gap suppression is observed as an increased resistance.

In conclusion, our observations put a new perspective on the feasibility of the superconducting spin switch. Close to \( T_c \), spin switch effects can be found with weak ferromagnets but the difficulty of obtaining highly transparent interfaces when using strongly disordered alloys may preclude full switching. Increasing the polarization, however, leads to a competing effect, namely the increased quenching of the superconductivity when the AP configuration reflects more spin-polarized quasiparticles back into the superconductor than the P-configuration. Note that both mechanisms are not mutually exclusive since one depends on the quasiparticles and the other on the Cooper pairs. We also have made clear the differences between CIP and CPP experiments. Finally, our observations should be of importance for the reproducibility and interpretation of data from S/F multilayers with strong magnets in general. In many reported experiments on \( T_c \)-variations, the magnetization state of the sample is undefined, and in particular re-entrant effects close to the critical thickness might be affected by the domain state of the sample.

We thank J. Santamaria, C. Bell and M. Hesselberth for useful discussions, and M. Hesselberth for help in the sample preparation. This work is part of the research program of the "Stichting voor Fundamenteel Onderzoek der Materie (FOM)", which is financially supported by NWO. The ESF-programs 'Pilshift' and 'Thixo' are acknowledged for providing an invaluable forum for discussing the preliminary results.

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