**Sc$_2$Ga$_2$CuO$_7$: A possible Quantum spin liquid near the percolation threshold**

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**Abstract**

Sc$_2$Ga$_2$CuO$_7$ (SGCO) crystallizes in a hexagonal structure (space group: $P6_3/mmc$), which can be seen as an alternating stacking of single and double triangular layers. Combining neutron, x-ray, and resonant x-ray diffraction we establish that the single triangular layers are mainly populated by non-magnetic Ga$^{2+}$ ions (85% Ga and 15% Cu), while the bi-layers have comparable population of Cu$^{2+}$ and Ga$^{3+}$ ions (43% Cu and 57% Ga). Our susceptibility measurements in the temperature range 1.8 - 400 K give no indication of any spin-freezing or magnetic long-range order (LRO). We infer an effective paramagnetic moment $\mu_{eff}=1.79 \pm 0.09 \mu_B$ and a Curie-Weiss temperature $\theta_{CW}$ of about $-44$ K, suggesting antiferromagnetic interactions between the Cu$^{2+}$ ($S=1/2$) ions. Low-temperature neutron powder diffraction data showed no evidence for LRO down to 1.5 K. In our specific heat data as well, no anomalies were found down to 0.35 K, in the field range 0-140 kOe. The magnetic specific heat, $C_m$, exhibits a broad maximum at around 2.5 K followed by a nearly power law $C_m \propto T^n$ behavior at lower temperatures, with $\alpha$ increasing from 0.3 to 1.9 as a function of field for fields up to 90 kOe and then remaining at 1.9 for fields up to 140 kOe. Our results point to a disordered ground state in SGCO.

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**Introduction:** Geometrically frustrated spin systems are full of surprises and continue to draw attention because of their tendency to host novel ground states [1]. In particular, triangular lattice based geometrically frustrated Mott insulators have gained wide interest because of their tendency to host novel ground states [1]. The unconventional nature of its elementary excitations, which result from a resonating valence bond (RVB) picture [2]. Anderson’s RVB proposal [2] pertained to the two dimensional (2D) edge-shared triangular lattice which, however, has a 120° ordered ground state [2,5]. This, in particular, triggered the search for the RVB state (also called the quantum spin liquid QSL state [1]) in geometries such as triangular (anisotropic), Kagomé (2D), hyperkagomé (three dimensional 3D), and pyrochlore (3D). Finally, a QSL state was realized notably in the 2D triangular system $\kappa$-(BEDT-TTF)$_2$Cu$_2$(CN)$_3$ [7], the Kagomé system ZnCu$_3$(OH)$_6$Cl$_2$ [8,12], and the hyperkagomé system Na$_4$Ir$_3$O$_7$ [13,14]. A frustrated geometry and a low value of spin ($S=1/2$), which enhances quantum fluctuations, help in stabilizing a QSL state [1]. However, even in the 6H-B and 3C phases of Ba$_3$NiSb$_2$O$_9$ (Ni$^{2+}$, $S=1$) [15], which have 2D edge-shared triangular and 3D edge-shared tetrahedral lattices, respectively, a QSL state has been suggested. Also, the realization of a QSL state for Ba$_4$IrTi$_2$O$_9$ [16], containing a diluted triangular lattice, and Ba$_3$YIr$_2$O$_9$ (high pressure cubic phase) [17], possibly suggests the importance of further neighbor interaction and/or deviations from the Heisenberg model. Recent experimental/theoretical results suggest that disorder might even drive the QSL state [18,19], the unconventional nature of its elementary excitations, which result from a chargeless sector of spin-1/2 fermions (commonly known as spinons), also drew interest from theorists and experimentalists [1,20,21]. Spinons are believed to form a Fermi surface in such Mott insulators and a non-vanishing Sommerfeld coefficient ($\gamma$) appears to be a generic feature for most of the previously discussed QSL [13,15,17].

We have been exploring a variety of spin systems with the objective of finding new QSL especially investigating triangular lattices. Herein, we introduce a Cu$^{2+}$ ($S=1/2$) based potential QSL system Sc$_2$Ga$_2$CuO$_7$ (SGCO). This system was first reported by Kimizuka et al. [22] and the structure type was identified as similar to Yb$_2$Fe$_3$O$_7$ [23]. However, no other data have been reported up to now. Here we report a thorough investigation of the structure and magnetic properties of SGCO by synchrotron x-ray diffraction (xrd), neutron diffraction (ND), susceptibility, and specific heat measurements. In the Yb$_2$Fe$_3$O$_7$ structure as applied to SGCO, the 4f sites

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Several Data were also of the Swiss Light Source (SLS), with the wavelength |

The powder xrd and Materials Sciences Beamline (see \textsuperscript{24}). On the other hand, the 4f sites (1/3, 2/3, 0.54) form triangular bi-planes (see Fig. 1 (c)) or rather double layers of triangular bi-pyramids when the oxygens are included. A three phase Rietveld refinement using FullProf \textsuperscript{27} was done on our synchrotron data to maintain clarity, not all atoms are shown). The 2b and 4f sites are shared by Cu and Ga (see legend and text). Since Cu and Ga have similar scattering lengths, both for x-rays and for neutrons, possible site sharing between the two is not easily addressed by analyzing the diffraction data. For instance, for x-rays (\(\lambda = 0.62 \, \text{\AA}\)) the atomic scattering factors for Ga and Cu are different by only 7%. However, if one tunes the x-ray energy near the K-absorption edge of Cu, the difference in scattering factors for Ga and Cu can be much more. In our case (\(\lambda = 1.38455 \, \text{\AA}\)), the scattering factors for Ga and Cu are different by 25%. Xrd pattern under these conditions is shown in Fig. 2. The obtained lattice constants are \(a = b = 3.3045(2) \, \text{\AA}, c = 28.1129(1) \, \text{\AA}\) (see Supplementary Material \textsuperscript{24} for details). We find that the 2b sites (0, 0, 1/4) contain 85% non-magnetic Ga\(^{3+}\) and 15% Cu\(^{2+}\). The 4f sites (1/3, 2/3, 0.54) contain 57% Ga\(^{3+}\) and 43% Cu\(^{2+}\) ions. As shown in Fig. 1 (b) the edge-shared single-triangular planes at the 2b sites (0, 0, 1/4), along with the neighboring oxygens atoms make single layers of corner-sharing triangular bi-pyramids (see Supplementary Material \textsuperscript{24}). On the other hand, the 4f sites (1/3, 2/3, 0.54) form triangular bi-planes (see Fig. 1 (c)) or rather double layers of triangular bi-pyramids when the oxygens are included. In a bi-plane (4f sites), a triangular layer is shifted with respect to the other such that the vertices of one layer are at the centroids of alternate triangles of the other layer. This is topologically equivalent to a honeycomb lattice with nearest (\(J1\)) and next-nearest (\(J2\)) neighbor coupling and based on exchange paths it is likely that \(J2 > J1\), see Fig. 1 (d).
The preference for Ga for the 2b sites is supported by bond-valence sum (BVS) calculations. The BVS for Ga$^{3+}$ is 2.91 on 2b sites and 2.69 on 4f sites, making 2b the preferred site for Ga$^{3+}$. Likewise the BVS for Cu is 2.54 on 2b sites and 2.35 on 4f sites, giving a slight preference for Cu$^{2+}$ to occupy 4f sites, in agreement with the results from the structural refinements.

The dc susceptibility ($\chi = \frac{M}{H}$) of SGCO varies in a Curie-Weiss manner $\chi = \chi_0 + C/(T - \theta_{CW})$ down to 100 K, where $\chi_0$ (-4.451 x 10^{-5} cm$^3$/mol), $C = 0.41$ cm$^3$/K/mol, and $\theta_{CW} = -44$ K denote the $T$-independent $\chi$, the Curie constant, and the Curie-Weiss temperature, respectively. The negative $\theta_{CW}$ suggests antiferromagnetic correlations and the inferred effective moment $\mu_{eff} \approx \sqrt{SC} = 1.79 \pm 0.09 \mu_B$ is close to that for an $S = 1/2$ moment. Measurements performed in a low field (25 Oe) and down to 2 K do not exhibit any anomaly or bifurcation between zero field cooled (ZFC) and field cooled (FC) data (inset of Fig. 3). AC $\chi$ measurements in the $T$-range 0.3-30 K do not show any anomaly or frequency dependence (see the Supplemental material). This suggests the absence of LRO or any spin freezing. It is worth mentioning that the ND measurements carried out at 1.5 K did not show any signature of magnetic Bragg peaks, setting an upper limit of 0.5 $\mu_B$ for the ordered moment in case of LRO. Further, $^{71}$Ga NMR shift data on SGCO indicate a leveling off of $\chi$ below about 50 K and the absence of LRO. While further evidence is needed to conclusively establish QSL behavior in SGCO, the essential features of $\chi(T)$ in SGCO are as in other QSL materials.

To estimate the fraction of paramagnetic spins in SGCO, we have made use of the data from magnetic isotherms $M(H)$ at low-$T$ (1.8 - 4.5 K) since the Cu at the 2b sites is expected to show paramagnetic Curie-like behavior while the contribution from the possibly correlated triangular bi-planes might be much lower. Our analysis (see the Supplemental material) of the $M(H)$ data is consistent with about 12% of paramagnetic Cu$^{2+}$ ($S = 1/2$) spins which is not far from the value of 15% Cu at the 2b sites obtained from resonant xrd data. In In$_2$Ga$_2$CuO$_7$, structurally similar to SGCO, Taetz et al. obtained approximately 10% of paramagnetic $S = 1/2$ impurities from a similar analysis. Given the site occupancies obtained from resonant xrd results on SGCO, the magnetic lattice (bi-layers from the atoms at 4f positions (1/3, 2/3, 0.54)) is more than 50% diluted. While this is beyond the percolation threshold of a 2D triangular lattice, for the bi-layer configuration (or the honeycomb lattice as mentioned before) here, the magnetic connectivity might still be maintained.

Based on our experiments, we suggest that SGCO appears to consist of (i) 10-15% of the 2b sites having $S = 1/2$ paramagnetic moments and (ii) a nearly equal mixture of $S = 1/2$ Cu$^{2+}$ and $S = 0$ Ga$^{3+}$ at the 4f sites which magnetically form the triangular bi-planes. While it could be that the site occupation is correlated, maintaining the magnetic connectivity in the triangular bi-layers, it is surprising that there is no spin-freezing even...
after large dilution. One may speculate that the proximity of the site occupation to a percolation threshold may promote long range spin singlets as in an algebraic spin liquid [1].

Heat capacity. To obtain further insight into the low-energy excitations of SGCO, we measured its specific heat \( C_p(T) \) (Fig. 1) in the \( T \)-range 0.35 - 300 K under various \( H \) (0 kOe - 140 kOe). We found no signature of LRO down to 0.35 K. However, below about 15 K the \( C_p(T) \) displays a dependence on \( H \) which is suggestive of a Schottky anomaly. This Schottky contribution \( C_{\text{Schottky}} \) most likely arises from the paramagnetic Cu\(^{2+} \) spins at the 2b sites. Note that our \( M(H) \) data analysis revealed a \( \sim 12\% \) contribution from paramagnetic \( S = 1/2 \) species. We then analyzed \( C_p(T) \) as a combination of \( C_{\text{Schottky}} \) and \( C_{\text{lattice}} \) (the lattice heat capacity) in addition to a magnetic heat capacity \( C_m \) (see Supplemental Material [24]). The \( C_m \) (which, we believe, comes from the correlated spins of the triangular bi-planes) is obtained after subtracting \( C_{\text{lattice}} \) and \( C_{\text{Schottky}} \) and is shown in Fig. 1(b).

![Figure 4: (Color online) (a) The \( C_p(T) \) data (log-log scale) for various \( H \) (up to 140 kOe) are shown. Inset depicts the appearance of a broad maximum in \( C_m \), in \( H = 0 \). (b) Power law fits (see text) of \( C_m \) are shown as solid lines. (c) Variation of exponent \( \alpha \) with \( H \) is shown. Note that, for \( H > 40 \) kOe the size of the error bars is equal to the size of the symbols.

As shown in the inset of Fig. 1(a), \( C_m(T) \) shows a broad maximum (perhaps not related to any phase transition) at around 2.5 K in \( H = 0 \). The appearance of a broad maximum is common for highly frustrated spin systems and has previously been noticed for NiGa\(_2\)S\(_4\) [30], Na\(_4\)Ir\(_3\)O\(_8\) [13], Ba\(_3\)CuSb\(_2\)O\(_8\) [31], and Ba\(_3\)NiSb\(_2\)O\(_9\) [15]. However, its position in \( T \) (in comparison to the strength of the exchange coupling) varies from material to material.

Finally the low-\( T \) \( C_m \) data \( (T < 1 \) K) were fitted to a power-law \( C_m = \gamma T^\alpha \) (where \( \gamma \) is a constant), to infer about the magnetic excitations (see Fig. 1(b)). We found that in SGCO, \( \alpha \) increases from 0.3 to 1.9 with increasing \( H \). The error bars for \( \alpha \) in Fig. 1(c) were arrived at by assuming a \( \pm 50\% \) uncertainty in \( C_{\text{Schottky}} \). The value of \( \alpha \) is robust and reliable for \( H > 40 \) kOe in the \( T \)-range 0.35-1 K, and attains a value 1.9(2) for \( H > 90 \) kOe. An exponent close to 2 might be an indication of a QSL and has previously been observed for the hyperkagomé compound Na\(_3\)Ir\(_3\)O\(_8\) [13] as also the 3C phase of Ba\(_3\)NiSb\(_2\)O\(_9\) [12]. In the case of a Kagomé structure, a theoretical approach (spin-1/2 spinons obeying the Dirac spectrum [32]) leads to a \( T^2 \) dependence of specific heat. However, there could be other explanations. Recently, randomness-induced QSL (also called random singlet) has been suggested for organic triangular salts [33] as also for inorganic ones such as herbertsmithite [34] with the Kagomé structure. In herbertsmithite, a Zn/Cu disorder might cause a random Jahn-Teller distortion of the [Cu(OH)\(_6\)]\(^{4-}\) giving rise to a random modification of the exchange coupling in the Kagomé layer. For the resulting gapless QSL state, a \( T \)-linear \( C_m \) is expected. In the SGCO case, a similar Ga/Cu disorder exists which might drive the QSL state. However, in SGCO (in zero field), the power-law exponent of \( C_m \) is 0.3. This amounts to a divergence of \( C_m/T \) at low-\( T \). Such a divergence in Pr\(_3\)Ir\(_2\)O\(_7\) (frustrated moments in a metallic background; different than SGCO) has been seen as an indication of proximity to a quantum critical point QCP [35]. For SGCO, with the increasing \( C_m(T) \) exponent (from 0.3 to 1.9) with field, one might speculate that the system moves from a QCP to deeper into the insulating QSL region with field. A cautionary note is that the inferred low-field \( C_m \) might be affected by possible interaction between the correlated spins and the orphan spins.

It is worth mentioning that in our \( C_m \) data at \( H = 140 \) kOe, \( \sim 34 \) mJ/mol K\(^2\). A non zero value of \( \gamma \), in general, indicates the presence of low-energy (likely) gapless excitations for these Mott insulators at low-\( T \). While the reason for the field-induced suppression of the \( C_m \) (suggesting a suppression of magnetic excitations) is not clear at this point, it could arise from a freezing out of any interaction between the orphan spins (Cu\(^{2+} \) in the 2b sites) and the triangular bi-layers. The estimated entropy change \( \Delta S \) (see Supplemental Material [24]) by integrating the \( C_m/T \) versus \( T \) data (in zero field) yields 1.2 J/mol K, which is nearly 20% of that expected for a \( S = 1/2 \) moment. This value will be even smaller when the data at the highest field are considered. The significantly smaller \( \Delta S \) is suggestive of a large residual entropy at low temperatures, further suggesting a QSL ground state in SGCO.

Summary: Using the results obtained from various experimental probes such as xrd, ND, dc/ac \( \chi(T) \), \( M(H) \), and \( C_p(T) \), we have explored the properties of SGCO. Our resonant xrd measurements carried out at the K-
absorption edge of Cu allow us to infer the site occupancies of Cu and Ga at the 2b sites (Cu:Ga = 0.15:0.85) and the 4f sites (Cu:Ga = 0.43:0.57). We suggest that the magnetic lattice can be viewed as a combination of (i) highly depleted triangular bi-layers (or honeycomb layers with nearest as also next-nearest neighbor interactions) giving rise to correlated behavior and (ii) about 15% Cu spins which are paramagnetic. Our ND data down to 1.5 K do not show any signature of LRO/spin freezing and are consistent with the dc/ac $\chi(T)$ measurements. Further, the $C_p(T)$ data indicate absence of any transition down to 0.35 K. They are also indicative of the presence of low-energy excitations and for high fields yield $C_m \propto T^2$ below 1 K. Taken together, the apparent lack of LRO or spin freezing and the existence of low-energy (likely gapless) excitations could be indicative of a QSL ground state. Often QSL candidates with site-disorder have been found to display spin freezing. It is therefore unusual to find the combined lack of order and spin-freezing in a system so heavily diluted as the triangular bi-layers are in this compound. This could be a result of the dilution approaching a percolation threshold and one may speculate if the unusual QSL behavior could arise close to such thresholds. There is also recent work [19] which proposes disorder driven bond randomness which can give rise to a random singlet QSL as has been suggested in herbertsmithite [22]. Further low-$T$ magnetization and local probe investigations such as $\mu$SR would be useful to explore the properties of SGCO in greater details.

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Supplemental material for Sc$_2$Ga$_2$CuO$_5$: A possible Quantum spin liquid near the percolation threshold

A. Sample preparation

Stoichiometric mixtures of pre-dried CuO (99.995% purity), Ga$_2$O$_3$(99.99% purity), and Sc$_2$O$_3$(99.9% purity) were ground homogeneously in an agate mortar and then pelletized. The pellet was placed in a high density alumina crucible and fired at 950°C for 24 hours in a Carbolite tubular furnace. The pellet was further reground, pelletized, and then fired two times at 1250°C for 24 hours each. Several such batches were prepared in this study. The final product was light green in color.

B. X-ray (synchrotron) diffraction and neutron diffraction analysis

Figure 5 depicts the x-ray diffraction pattern obtained for SGCO using synchrotron data with $\lambda = 0.6204$ Å. In the ideal case (no site sharing), the Cu and Ga atoms are expected to form single and double triangular lattices comprising of atoms at 2b (0, 0, 1/4) and 4f (1/3, 2/3, z) sites, respectively. However, because of the similar ionic radii of Cu$^{2+}$ and Ga$^{3+}$ ions, a distribution of both might be expected at the 2b (0, 0, 1/4) and the 4f (1/3, 2/3, z) sites. The atomic coordinates refined from the x-ray (synchrotron: SLS) diffraction data measured with $\lambda = 0.6204$ Å at 300 K for SGCO are given in Table II. Since Cu (Z = 29) and Ga (Z = 31) have nearly the same atomic scattering factors it is difficult to estimate with high confidence, their distribution with the present x-ray (synchrotron: $\lambda = 0.6204$ Å) analysis. If the Ga:Cu ratio were left as a free parameter, the best refinement of the synchrotron (SLS) data was obtained with the site 2b (0, 0, 1/4) mostly occupied by Ga$^{3+}$ and the 4f (1/3, 2/3, z) site having both Cu$^{2+}$ and Ga$^{3+}$ ions, which is quite opposite to the isostructural compound Yb$_2$Fe$_3$O$_7$ where the divalent Fe$^{2+}$ and the trivalent Fe$^{3+}$ ions sit at 2b (0, 0, 1/4) and 4f (1/3, 2/3, z) sites having both Cu$^{2+}$ and Ga$^{3+}$ ions, which is quite opposite to the isostructural compound Yb$_2$Fe$_3$O$_7$ where the divalent Fe$^{2+}$ and the trivalent Fe$^{3+}$ ions sit at 2b (0, 0, 1/4) and 4f (1/3, 2/3, z) sites, respectively. The refinement quality factors described by $R_p$, $R_{exp}$, $R_{wp}$ and $\chi^2$ are obtained to be 3.56%, 0.34%, 3.50% and 108, respectively. The local environments of Cu$^{2+}$ ions at 2b/4f site can be best viewed as single layers of corner-sharing triangular bi-pyramids/double layers of triangular bi-pyramids, as seen in Fig. 6. As far as the magnetic lattice is concerned, the latter amounts to a (heavily depleted) honeycomb lattice with the presence of both nearest and next-nearest neighbor interactions.

We then collected x-ray diffraction data near the K absorption edge of Cu to get a more reliable distribution of Cu and Ga atoms at 2b and 4f sites, as the contrast between the atomic scattering factors of Ga and Cu is higher at this wavelength. From a refinement of our resonant x-ray data with $\lambda = 1.38455$ Å (refined positions and occupancies are in Table III), we obtain that the 2b sites (0, 0, 1/4) are mostly occupied with non-magnetic Ga$^{3+}$ ions (85% Ga$^{3+}$ and 15% Cu$^{2+}$), whereas the 4f sites (1/3, 2/3, z), which form a network of double layers of triangular bi-pyramids (or honeycomb layers as mentioned above), is populated by Ga$^{3+}$ and Cu$^{2+}$ ions with 57% Ga$^{3+}$ and 43% Cu$^{2+}$, see Fig. 6. The goodness of fit parameters $R_p$, $R_{exp}$, $R_{wp}$ and $\chi^2$ are found to be 2.67%, 2.94%, 0.34% and 75, respectively.

![XRD (SLS)-300 K](image)

Figure 5: (Color online) X-ray (synchrotron) diffraction data for SGCO recorded at 300 K with wavelength 0.6204 Å. Iobs and Ical represent the experimental and Rietveld refined intensities, respectively. The vertical blue, pink and green bars depict the Bragg positions for Sc$_2$Ga$_2$CuO$_5$, Sc$_2$O$_3$ and CuGa$_2$O$_4$, respectively, whereas flat cyan line shows the intensity difference Iobs- Ical.

Table I: Atomic positions obtained after Rietveld refinement of the x-ray (synchrotron:SLS) data at 300 K with $\lambda = 0.6204$ Å.

| Atom | Site | x   | y   | z   | occ |
|------|------|-----|-----|-----|-----|
| Sc   | 4f   | 1/3 | 2/3 | 0.14877(2) | 1   |
| Cu   | 2b   | 0.0 | 0.0 | 1/4 | 0.252(2) |
| Ga   | 2b   | 0.0 | 0.0 | 1/4 | 0.744(2) |
| Ga   | 4f   | 1/3 | 2/3 | 0.54478(1) | 0.312(3) |
| Cu   | 4f   | 1/3 | 2/3 | 0.54478(1) | 0.69(3) |
| O1   | 4f   | 0.0000 | 0.0000 | 0.18007(6) | 1   |
| O2   | 4f   | 1/3 | 2/3 | 0.613(30) | 6   |
| O3   | 4f   | 1/3 | 2/3 | 0.03325(7) | 1   |
| O4   | 2c   | 1/3 | 2/3 | 1/4 | 1   |

arXiv:1506.03576v2.
[35] Y. Tokiwa, J. J. Ishikawa, S. Nakatsuji, and P. Gegenwart, Nature Materials 13, 359 (2014).
Figure 6: (Color online). (a) Local arrangement of atoms in a unit cell of SGCO is shown. Blue and magenta color balls each correspond to a fractional occupation by the Cu and Ga atoms as mentioned in the legend. Figures 1 (b) and 1 (c) represent the Cu/Ga environments at the 2b (corner-sharing triangular bipyramids) and the 4f (double layers of triangular bipyramids) sites, respectively.

Table II: Crystal structure parameters of SGCO as refined from synchrotron data ($\lambda = 1.38455$ Å) at 300 K with space group $P6_3/mmc$. The occupancies of Cu and Ga at both sites have been constrained to ensure the ratio of total amounts of Cu and Ga to be 1:2.

| Atom | Wyckoff position | x   | y   | z     | occupancy |
|------|------------------|-----|-----|-------|-----------|
| Sc   | 4f               | 1/3 | 2/3 | 0.14865(2) | 1         |
| Cu1/Ga1 | 2b           | 0   | 0   | 1/4   | 0.66      |
| Cu2/Ga2 | 4f           | 1/3 | 2/3 | 0.54480(1) | 0.16      |
| O1   | 4f               | 0   | 0   | 0.1763(4)   | 1         |
| O2   | 4f               | 1/3 | 2/3 | 0.6118(7)   | 1         |
| O3   | 4f               | 1/3 | 2/3 | 0.03371(7)  | 1         |
| O4   | 2c               | 1/3 | 2/3 | 1/4   | 1         |

Similarly, the neutron diffraction measurements performed at 300 K, shown in Fig. 7, also could not reliably infer the amount of site sharing between Cu$^{2+}$ and Ga$^{3+}$ ions, because of a poor contrast between their scattering lengths, 7.718 and 7.288 fm, respectively [1]. Consequently, based on our neutron diffraction data alone we could not have made a definitive statement about the site occupancy by Ga and Cu at the 2b and 4f sites. The obtained lattice constants from neutron diffraction data ($\lambda = 1.49$ Å) are $a = b = 3.30459(4)$ Å and $c = 28.1183(4)$ Å, in good agreement with those obtained from x-ray (synchrotron) data. The residual refinement factors $R_p$, $R_{exp}$, $R_{wp}$ and $\chi^2$ for neutron refined powder diffraction pattern for 300K data yield 3.61%, 3.00%, 4.73% and 2.49, respectively. The refined atomic coordinates for

Figure 7: (Color online) Neutron diffraction data measured for SGCO with $\lambda = 1.49$ Å at 300 K are shown. The open black circles and the red solid line represent the experimental and refined data. The vertical blue and pink lines indicate the Bragg positions for SGCO and Sc$_2$O$_3$, while flat cyan line shows the difference plot between measured and refined intensities.

Table III: Refined positional parameters and the occupancies for SGCO at 300 K under the space group $P6_3/mmc$ using the neutron diffraction data with $\lambda = 1.49$ Å.

| Atom | Wyckoff position | x   | y   | z     | occupancy |
|------|------------------|-----|-----|-------|-----------|
| Sc   | 4f               | 1/3 | 2/3 | 0.14861(5) | 0.93      |
| Cu   | 2b               | 0   | 0   | 1/4   | 0.66      |
| Ga   | 2b               | 0   | 0   | 1/4   | 0.34      |
| Cu   | 4f               | 1/3 | 2/3 | 0.54447(5) | 0.16      |
| Ga   | 4f               | 1/3 | 2/3 | 0.54447(5) | 0.84      |
| O1   | 4f               | 0   | 0   | 0.18087(7) | 1         |
| O2   | 4f               | 1/3 | 2/3 | 0.61201(6) | 1         |
| O3   | 4f               | 1/3 | 2/3 | 0.03320(10) | 1         |
| O4   | 2c               | 1/3 | 2/3 | 1/4   | 1         |
SGCO at 300 K are listed in the Table III. As it is evident from the refinement result that the neutron diffraction suggests a distribution of Cu and Ga quite opposite to what was determined from resonant x-ray data. The neutron diffraction data collected at 1.5 K (Fig. S) do not show the appearance of any new peak and hence rule out the presence of magnetic ordering.

Figure 8: (Color online) Neutron diffraction data for SGCO, collected at 1.5 K, with its Rietveld refinement is shown. Vertical blue and pink bars depict the Bragg positions for SGCO (main phase) and Sc$_2$O$_3$ (impurity phase \( \sim 1.2\% \)).

Figure 9: (Color online) The magnetization isotherms measured in the field range (0-70 kOe) and in the temperature range 1.8-4.5 K are shown. The dotted broken lines are fits to an equation described in the text. The panels (a), (b), (c), and (d) illustrate these fits for various values of the spectroscopic splitting factor \( g = 2.0, 2.1, 2.2, \) and 2.3, respectively.

Table IV: Results of various fits to the magnetization isotherm data with the equation described in text.

| \( S_{\text{imp}} \) | \( g_{\text{imp}} \) | \( \chi(\text{cm}^3/\text{mol}) \) | \( \theta_{\text{imp}}(\text{K}) \) | \( f_{\text{imp}}(\%) \) |
|-----------------|-----------------|-----------------|-----------------|-----------------|
| 1/2             | 2.0             | 0.0051          | -0.22           | 11.60           |
| 1/2             | 2.1             | 0.0047          | -0.42           | 11.58           |
| 1/2             | 2.2             | 0.0043          | -0.63           | 11.59           |
| 1/2             | 2.3             | 0.0038          | -0.85           | 11.61           |

C. Magnetization isotherm analysis

As it is evident from Fig. 9 below about 5 K magnetization isotherms start to show a curvature, which becomes prominent at lower temperatures. Measured magnetic isotherms in the temperature range 1.8-4.5 K, where \( \mu_B H \) is no longer much smaller than \( k_B T \), were analyzed within the framework of the following equation: \( M(H,T) = \chi(T)H + N_A\mu_B f_{\text{imp}} S_{\text{imp}} g_{\text{imp}} B_S(g_{\text{imp}} S_{\text{imp}} H \mu_B/k_B(T - \theta_{\text{imp}})) \), where \( \chi \) represents the intrinsic susceptibility of the triangular planes and is assumed constant in the 1.8-4.5 K interval, \( f_{\text{imp}}, g_{\text{imp}}, S_{\text{imp}}, \) and \( \theta_{\text{imp}} \) denote the impurity/disorder concentration, the Lande \( g \)-factor, spin, and correlation temperature for impurity spins. The symbols \( H, N_A, \mu_B, B_S \) refer to the magnetic field, the Avogadro number, the Boltzmann constant, the Bohr magneton, and the Brillouin function, respectively.

To start with, there are five-parameters to be varied and a simultaneous fitting of all these end up in nonphysical values. So we first fixed \( S_{\text{imp}} \) and \( g_{\text{imp}} \) to be 1/2 and 2, respectively and thus reduced our variables from five to three, i.e., \( \chi, f_{\text{imp}} \) and \( \theta_{\text{imp}} \). To increase the reliability of our fitting parameters we simultaneously fitted all magnetic isotherms, measured in the \( T \)-range 1.8 - 4.5 K. The outcomes of fittings for various assumed values of \( g \) are summarized in Table IV. As is evident from the result, the value of \( \theta_{\text{imp}} \) is nearly zero, which in turn suggests the absence of magnetic correlations among impurity spins so one can treat them as purely paramagnetic. The obtained free-spin concentration is consistent with about 12% antisite disorder.

D. A.C. susceptibility

AC susceptibility data were measured in the temperature range 0.3 - 7 K and 2- 30 K on a home-made susceptometer and a Quantum Design SQUID VSM machine, respectively. The data were recorded keeping the dc field zero and the ac pulse of amplitude 5 Oe. With the homemade susceptometer we measured the data for three different frequencies of 545.7 Hz, 311 Hz, and 155 Hz. In the temperature range 2-30 K on Quantum design SQUID VSM machine we measured the data for five different frequencies of 11 Hz, 110 Hz, 155 Hz, 311 Hz, and 546 Hz.
used a combination of Debye and Einstein terms.

In the absence of a suitable non-magnetic analog for SGCO, we cannot estimate the magnetic spins (Cu$^{2+}$) in the 4f planes (bi-planes). We now attempt to estimate the contribution of the magnetic spins residing in the 4f planes.

The specific heat of the system can be written as:

$$C(T) = C_{lattice} + C_{Schottky} + C_m$$

where, $C_{lattice}$ is the lattice specific heat and $C_{Schottky}$ originates from isolated paramagnetic spins (Cu$^{2+}$ spins in this case) forming a two-level system. We ascribe the remaining contribution $C_m$ to antiferromagnetically interacting Cu$^{2+}$ spins residing in the 4f planes (bi-planes). We now attempt to estimate $C_{lattice}$ and $C_{Schottky}$ so that the “intrinsic” contribution $C_m$ can be inferred.

To estimate the lattice contribution ($C_{lattice}$) in the absence of a suitable non-magnetic analog for SGCO, we used a combination of Debye and Einstein terms, as written in Eqn. (1) given below.

$$C_{lattice}(T) = C_{Debye} + C_{Einstein}$$

where, $C_{Debye} = C_D \left[ 9k_B \left( \frac{\pi D}{x_D} \right)^3 \sum \frac{x^4 e^{-x}}{(e^x - 1)^2} dx \right]$ and $C_{Einstein} = \sum C_{E_i} \left[ 3R \left( \frac{\theta_E}{T} \right)^2 \frac{\exp\left( \frac{\theta_E}{T} \right)}{\exp\left( \frac{\theta_E}{T} \right) - 1} \right]$

One formula unit of SGCO has twelve atoms and this in turn offers one acoustic and eleven optical modes of atomic vibrations, in each crystallographic direction. The first term (Debye integral) of the Eqn. (1) takes care of the three acoustic modes of vibration whereas the second term accounts for the contribution of thirty three optical modes of phonons. This amounts to the contribution of one acoustic branch of the phonon spectrum) dominates as $T^3$ and Einstein weightage to the lattice specific heat falls rapidly as $\exp(-\theta_D/T)$.

Figure 10: Variation of the real part of the ac susceptibility as a function of temperature (0.3-30 K) for several frequencies is shown.

### E. Magnetic specific heat analysis

The $C_p(T)$ data, specific heat at constant pressure, data obtained in zero magnetic field in the temperature range 0.35-300 K are shown in Fig. 11. The total specific heat of the system can be written as: $C_p(T) = C_{lattice} + C_{Schottky} + C_m$. Here, $C_{lattice}$ is the lattice specific heat and $C_{Schottky}$ originates from isolated para-magnetic spins (Cu$^{2+}$ spins in this case) forming a two-level system. We ascribe the remaining contribution $C_m$ to antiferromagnetically interacting Cu$^{2+}$ spins residing in the 4f planes (bi-planes). We now attempt to estimate $C_{lattice}$ and $C_{Schottky}$ so that the “intrinsic” contribution $C_m$ can be inferred.

To estimate the lattice contribution ($C_{lattice}$) in the absence of a suitable non-magnetic analog for SGCO, we used a combination of Debye and Einstein terms, as written in Eqn. (1) given below.

$$C_{lattice}(T) = C_{Debye} + C_{Einstein}$$

The first term (Debye integral) of the Eqn. (1) takes care of the three acoustic modes of vibration whereas the second term accounts for the contribution of thirty three optical modes of phonons. This amounts to the contribution of one acoustic branch of the phonon spectrum) dominates as $T^3$ and Einstein weightage to the lattice specific heat falls rapidly as $\exp(-\theta_D/T)$.

Figure 11: (Color online) Heat capacity at zero field is shown with black open circles. Cyan and red solid lines represent initial fit (20-90 K) to a combination of Debye and Einstein model and its extrapolation in the entire temperature range (0.35-300 K), respectively. Inset (a) highlights the low-temperature portion of the $C_p(T)$ data and inset (b) illustrates the entropy change at zero field.

Next, we will estimate the Schottky contribution $C_{Schottky}$ as follows. Based on magnetization isotherm analysis, we modeled the Schottky contribution ($C_{Schottky}$) by fixing the amount of free spins to 12%. We
used the two-level Schottky formula, described by Eqn. 2 below, with \( f = 12\% \) contribution from non-interacting Cu-spins having a \( g \)-value of 2.

\[
C_{\text{Schottky}} = f \left[ R \left( \frac{\Delta}{k_B T} \right)^2 \frac{g_0}{g_1} \frac{\exp \left( \frac{\Delta}{k_B T} \right)}{1 + \left( \frac{g_0}{g_1} \right) \exp \left( \frac{\Delta}{k_B T} \right)} \right]^{2}
\]  

(2)

Here, \( \Delta \) is the energy separation under the influence of an external magnetic field \( H \), and \( R, k_B, g_0 \) and \( g_1 \) are the universal gas constant, the Boltzmann constant, and the degeneracies of the two level system, respectively.

The 'intrinsic' magnetic specific heat \( C_m \) is thus obtained by subtracting \( C_{\text{lattice}} + C_{\text{Schottky}} \) from \( C_p(T) \).

Figure 12 illustrates a plot of individual contributions of Schottky (for \( g = 2 \) and 12% free spin contribution), lattice and magnetic specific heats to the total specific heat data. It is worth mentioning that below about 1 K the lattice heat capacity as also the Schottky contribution in fields beyond 90 kOe is utterly negligible in the total specific heat and the determination of the magnetic specific heat \( (C_m) \) below 1 K is free from any uncertainty.

Further, analyzing the magnetic heat capacity data assuming a \( \pm 50\% \) uncertainty in the Schottky part does not change significantly the values of exponents for fields \( H > 40 \) kOe, and the exponents are found to be robust for \( H > 40 \) kOe (see main text and Fig. 4 (c)). Note that Fig. 13 shows the lattice removed \( C_p(T) \) data, a sum of both the magnetic and Schottky specific heat, for SGCO at different magnetic fields.

Figure 13: (Color online) A sum of magnetic specific heat \( C_m \) and Schottky specific heat \( C_{\text{Schottky}} \) is shown for fields 0 - 140 kOe as a function of temperature.