Nickel-rich outflows from accretion discs formed by the accretion-induced collapse of white dwarfs

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ABSTRACT
A white dwarf (WD) approaching the Chandrasekhar mass may in several circumstances undergo accretion-induced collapse (AIC) to a neutron star (NS) before a thermonuclear explosion ensues. It has generally been assumed that such an AIC does not produce a detectable supernova (SN). If, however, the progenitor WD is rapidly rotating (as may be expected due to its prior accretion), a centrifugally supported disc forms around the NS upon collapse. We calculate the subsequent evolution of this accretion disc and its nuclear composition using time-dependent height-integrated simulations with initial conditions taken from the AIC calculations of Dessart and collaborators. Soon after its formation, the disc is cooled by neutrinos and its composition is driven neutron rich (electron fraction \(Y_e \sim 0.1\)) by electron captures. However, as the disc viscously spreads, it is irradiated by neutrinos from the central proto-NS, which dramatically alters its neutron-to-proton ratio. We find that electron neutrino captures increase \(Y_e\) to \(\sim 0.5\) by the time that weak interactions in the disc freeze out. Because the disc becomes radiatively inefficient and begins forming \(\alpha\)-particles soon after freeze out, powerful winds blow away most of the disc’s remaining mass. These \(Y_e \sim 0.5\) outflows synthesize up to a few times \(10^{-2} M_\odot\) in \(^{56}\text{Ni}\). As a result, AIC may be accompanied by a radioactively powered SN-like transient that peaks on a time-scale of \(\sim 1\) d. Since few intermediate mass elements are likely synthesized, these nickel-rich explosions should be spectroscopically distinct from other SNe. The time-scale, velocity and composition of the AIC transient can be modified if the disc wind sweeps up a \(\sim 0.1 M_\odot\) remnant disc created by a WD–WD merger; such an ‘enshrouded’ AIC may account for sub-luminous, sub-Chandrasekhar Type I SNe. Optical transient surveys such as the Panoramic Survey Telescope and Rapid Response System and the Palomar Transient Factory should detect a few AIC transients per year if their true rate is \(\sim 10^{-2}\) of the type Ia rate, and Large Synoptic Survey Telescope should detect several hundred per year. High cadence observations (\(\lesssim 1\) d) are optimal for the detection and followup of AIC.

Key words: accretion, accretion discs – neutrinos – nuclear reactions, nucleosynthesis, abundances – stars: neutron – supernovae: general.

1 INTRODUCTION
It is generally thought that when an accreting C–O white dwarf (WD) approaches the Chandrasekhar mass \(M_{Ch}\), the carbon ignites in its core, causing the WD to explode and produce a Type Ia supernova (SN). If, however, the WD’s composition is O–Ne instead of C–O, the outcome may be qualitatively different: electron captures on nuclei will cause the WD to collapse to a neutron star (NS) before an explosion can ensue (Miyaji et al. 1980; Canal, Isern & Labay 1990; Gütiérrez, Canal & García-Berro 2005; Poelarends et al. 2008). Such an accretion-induced collapse (AIC) may also occur for C–O WDs with large initial masses \((M \sim 1.1 M_\odot)\) in systems with mass-transfer rates \(\dot{M} \gtrsim 10^{-8} M_\odot\) yr\(^{-1}\) (Nomoto & Kondo 1991) or following the merger of two WDs in a binary if their total mass exceeds \(M_{Ch}\) (e.g. Mochkovitch & Livio 1989, 1990; Yoon, Podsiadlowski & Rosswog 2007). Indeed, spectroscopic surveys of WD–WD binaries (e.g. the SN Ia progenitor survey, SPY; Napiwotzki et al. 2004) have discovered at least one system that will merge within a few Hubble times with a total mass tantalizingly close to \(M_{Ch}\) (Napiwotzki et al. 2002).

Despite its likely occurrence in nature, AIC has not yet been observationally identified, probably because it is less common and produces less \(^{56}\text{Ni}\) than a normal SN (e.g. Woosley & Baron 1992). Identifying AIC or constraining its rate would, however, provide unique insights into the paths of degenerate binary evolution, the
true progenitors of Type Ia SNe (e.g. Yungelson & Livio 1998), and the formation channels of globular cluster NSs (Grindlay 1987), low-mass X-ray binaries (van den Heuvel 1984; Michel 1987), and millisecond pulsars (Grindlay & Bailyn 1988; Bailyn & Grindlay 1990). Identifying the optical signature of AIC is also important because these events may be a source of strong gravitational wave (GW) emission (Fryer, Holz & Hughes 2002; see Ott 2009 and references therein).

In this Letter, we show how AIC may itself produce a SN-like transient which, although less luminous and shorter lived than a typical SN, may not be detectable with upcoming optical transient surveys. Although little $^{56}\text{Ni}$ is ejected by the newly formed proto-NS itself, we show that powerful winds from the accretion discs formed during AIC produce up to a few times $10^{-2}\text{M}_\odot$ in $^{56}\text{Ni}$.

### 1.1 Nickel-rich winds from AIC discs

Fig. 1 summarizes the stages of AIC. Due to their preceding accretion, WDs that undergo AIC are probably rotating rapidly prior to collapse (Fig. 1A). As a result, a substantial fraction of the WD mass must be ejected into a disc during the collapse in order to conserve total angular momentum (Shapiro & Lightman 1976; Michel 1987). Dessart et al. (2006, 2007, hereafter D06, D07) perform two-dimensional axisymmetric AIC calculations, in which they find that a quasi-Keplerian accretion disc with mass $M_{A,0} \sim 0.1-0.5\text{M}_\odot$ forms around the newly formed proto-NS.

![Figure 1](https://academic.oup.com/mnras/article-abstract/396/3/1659/1746938)

**Figure 1.** Stages of the AIC of a WD. (A) WD accretes up to near the Chandrasekhar mass ($M_{\text{Ch}}$) from either a binary companion or following the disruption of another WD in a merger. Due to the accretion of angular momentum, the WD is rapidly rotating prior to collapse. (B) Electron captures cause the WD to collapse to a proto-NS surrounded by a compact accretion disc with a mass $\sim 0.1\text{M}_\odot$ and a size $\sim 100\text{km}$. Electron captures drive the composition neutron rich (i.e. to an electron fraction $Y_e \sim 0.1$). As matter begins to accrete onto the proto-NS, the bulk of the disc spreads to larger radii. (C) As the disc spreads, the mid-plane is irradiated by neutrinos from the central proto-NS. This irradiation is absorbed by neutrons, raising the electron fraction to $Y_e \sim 0.5$. (D) The disc's electron fraction freezes out. Soon thereafter, powerful outflows blow the disc apart. As the ejecta expands adiabatically, material with $Y_e \gtrsim 0.5$ is efficiently synthesized into $^{56}\text{Ni}$, the subsequent decay of which powers a SN-like optical transient.

In Metzger, Piro & Quataert (2008a, 2009 hereafter MPQ08, MPQ09), we studied the time-dependent evolution of accretion discs with similar initial properties that are formed from black hole (BH)–NS or NS–NS mergers. Many of these results also apply to proto-NS accretion following AIC. Although the disc is initially concentrated just outside the proto-NS surface, the bulk of the disc's mass must spread to larger radii as matter accretes in order to conserve angular momentum. Early in their evolution, such 'hyperaccreting' discs are neutrino cooled and geometrically thin, and their composition is driven neutron-rich by an equilibrium between electron and positron captures under degenerate conditions (e.g. Pruet, Woosley & Hoffman 2003; fig. 1B): the electron fraction $Y_e \sim 0.1$, where $Y_e = n_e/(n_p + n_0)$ and $n_p/n_0$ is the proto-neutron number density. As the disc viscously spreads and its temperature decreases, however, weak interactions eventually become slow compared to the evolution time-scale of the disc and the value of $Y_e$ freezes out.

Soon following freeze out, neutrino cooling becomes inefficient and the disc becomes advective and geometrically thick (MPQ08, MPQ09). Since advective discs are only marginally bound, powerful winds likely begin to unbind most of the disc's remaining mass (Blandford & Begelman 1999), aided by the nuclear energy released as $\alpha$-particles begin to form (Lee & Ramirez-Ruiz 2007; MPQ08, MPQ09). Because $\sim 20\%$ of the initial disc mass $M_{A,0}$ still remains when the disc becomes advective (or when $\alpha$-particles form), a robust consequence of disc formation during AIC is the ejection of up to a few times $10^{-2}\text{M}_\odot$. As the ejecta are hot and dense, heavy isotopes are synthesized as they expand away from the mid-plane and cool. Because weak interactions are already slow by this point, heavy isotopes are produced depends on $Y_e$ in the disc at freeze out.

In the case of BH accretion following BH–NS or NS–NS mergers, the disc freezes out neutron rich with an electron fraction $Y_e \sim 0.2-0.4$ (MPQ08; MPQ09). A crucial difference in the case of AIC, however, is the presence of the central proto-NS, which radiates a substantial flux of electron neutrinos as it deleptonizes during the first few seconds following core bounce (e.g. Burrows & Lattimer 1986). As $\nu_e$ flux irradiates the disc, which acts to raise $Y_e$ via absorptions on free neutrons ($\nu_e + n \rightarrow e^- + p$; Fig. 1C).

In Section 2, we present calculations of the evolution of the accretion discs formed from AIC using time-dependent height-integrated simulations with initial conditions and neutrino irradiation taken from the AIC simulations of D06. We show that due to $\nu_e$ absorptions a substantial fraction of the disc freezes out with $Y_e \gtrsim 0.5$. As a result, late time outflows from the disc primarily synthesize $^{56}\text{Ni}$ (Fig. 1D). We discuss the observational signature of these Ni-rich outflows in Section 3, and we conclude in Section 4 by evaluating the prospects for detecting AIC with upcoming optical transient surveys and as electromagnetic counterparts to GW sources.

### 2 AIC ACCRETION DISC MODEL

#### 2.1 Methodology

The equations and assumptions employed here closely follow those in section 3.1 of MPQ09 for NS–NS or BH–NS mergers. To provide a brief summary, we evolve the surface density $\Sigma$, temperature $T$ and electron fraction $Y_e$ as a function of radius $r$ and time $t$ using the 2N-RK3 scheme described in Brandenburg (2003). We employ an '\(\alpha\)' prescription for the viscosity (Shakura & Sunyaev 1973), with the accretion stress proportional to the gas pressure. Our calculation includes all of the relevant neutrino cooling processes for both...
optically thin and optically thick discs (e.g. Di Matteo, Perna & Narayan 2002). We take the initial surface density $\Sigma(r, t = 0)$ from the Keplerian disc that forms soon after core bounce in the AIC calculations of D06. We focus on the 1.46 $M_\odot$ WD model of D06, which collapses to form a $\approx 0.1$ $M_\odot$ disc that extends from the NS surface at $\approx 30$ km out to several NS radii (see fig. 5 of D06).

We evolve $Y_e$ including both local $e^-/e^+$ pair captures ($e^- + p \rightarrow v_e + n$ and $e^+ + n \rightarrow v_\bar{e} + p$) and non-local neutrino absorptions ($v_e + n \rightarrow e^- + p$ and $v_\bar{e} + p \rightarrow e^+ + n$) due to irradiation from the central proto-NS. We approximate the neutrino flux from the proto-NS as arising from a ‘light bulb’ at the origin because the neutrino capture rates only become important relative to the pair capture rates when the disc has spread to a radius which is much larger than that of the central NS. We neglect the disc’s neutrino luminosity because it is small compared to the proto-NS luminosity. We assume that the disc’s initial electron fraction is $Y_e(r, t = 0) = 0.5$, i.e. that of the WD prior to collapse.

The neutrino absorption rates are proportional to $L_{\nu_e}(\epsilon_{\nu_e})$, where $L_e$ and $\langle \epsilon_e \rangle$ are the neutrino luminosities and mean energies, respectively, from the proto-NS. During the first $t \approx 1$ s following core bounce, we use $L_{\nu_e}, L_{\nu_\bar{e}}, \langle \epsilon_{\nu_e} \rangle$ and $\langle \epsilon_{\nu_\bar{e}} \rangle$ from D06 (specifically, their values in the equator; see their fig. 9). The $L_{\nu_e}$ flux peaks at a very large value ($\gtrsim 10^{52}$ erg s$^{-1}$) at early times, after the shock created at core bounce ‘breaks out’ of the proto-NS’s neutrinosphere. The flux at later times ($t > 0.1$ s) is dominated by the cooling proto-NS, with $L_{\nu_e}$ remaining $\approx L_{\nu_\bar{e}}$ as the proto-NS deleptonizes (Burrows & Lattimer 1986). Note that because $L_{\nu_e}(\epsilon_{\nu_e}) > L_{\nu_\bar{e}}(\epsilon_{\nu_\bar{e}})$ during most of the first second following core bounce, the rate of $v_e$ captures exceeds the rate of $v_\bar{e}$ captures and the net effect of neutrino absorptions is to drive $Y_e$ to a value $\gtrsim 0.5$ (Qian et al. 1993). This result is robust since shock break out and deleptonization are generic features of NS formation (e.g. Burrows & Mazurek 1983); indeed, we find similar results using luminosities and mean energies from the rotating core-collapse simulations of Thompson, Quataert & Burrows (2005).

When necessary, for $t \gtrsim 1$ s we take $L_e \propto t^{-1}$, motivated by the calculations of Pons et al. (1999), and assume $\langle \epsilon_e \rangle \propto L_e^{4/3} \propto t^{-1/3}$, as appropriate for a blackbody with a fixed radius. Although the neutrino luminosities and spectra at such late times are somewhat uncertain, we find that for $\alpha \gtrsim 0.01$ the disc evolves sufficiently rapidly that our results are relatively insensitive to the precise values of $L_e$ and $\langle \epsilon_e \rangle$ at late times.

### 2.2 Results

The viscous and thermal evolution of the accretion discs formed from AIC are very similar to those presented in MPQ09 for NS–NS/NS–BH mergers, so we refer the reader to this work for a detailed discussion of the overall evolution of the disc. The major qualitative difference between AIC and NS–NS/NS–BH merger discs is the evolution of the electron fraction, so we focus on our results for $Y_e$ below.

Fig. 2 shows histograms of mass as a function of electron fraction $M(Y_e)$ for our fiducial model that employs an initial density profile and neutrino irradiation from the 1.46 $M_\odot$ WD model of D06. The solid and dashed lines show calculations performed with viscosity $\alpha = 0.3$ and 0.03, respectively. The top panel shows $M(Y_e)$ when 90 per cent of the disc’s mass has become advective, which corresponds to a time $t = 0.2(7)$ s for $\alpha = 0.3(0.03)$. Annuli in the disc are considered advective once the neutrino cooling rate becomes $\gtrsim 1/2$ of the viscous heating rate. The bottom panel shows $M(Y_e)$ when 90 per cent of the disc’s mass has fused into $\alpha$-particles, which corresponds to a time $t = (0.6)7$ s for $\alpha = 0.3(0.03)$. Although powerful outflows likely begin when the disc becomes advective (Blandford & Begelman 1999), thermonuclear disruption of the disc is assured once $\alpha$-particles form. This occurs because (1) the nuclear reactions that create (and destroy) deuterium (the building blocks of $\alpha$-particles) are still very fast compared to the evolution time-scale of the disc [i.e. the disc is still in nuclear statistical equilibrium (NSE)] when the temperature drops sufficiently that the NSE $\alpha$-particle mass fraction becomes approximately unity (e.g. Chen & Beloborodov 2007, equation 13) and (2) the energy released by forming $\alpha$-particles ($\sim 7$ MeV per nucleon) exceeds the disc’s gravitational binding energy (see MPQ08, MPQ09). Thus, the distributions in Fig. 2 roughly bracket the distribution of $Y_e$ in the outflows that carry away the remaining mass in the disc at late times.

Fig. 2 shows that $\approx 0.02–0.03$ M$_\odot$ (or $\approx 20–30$ per cent of the initial disc mass) remains when the disc is largely advective. A comparable amount of mass remains when the disc forms $\alpha$-particles because this occurs at approximately the same time (MPQ08; MPQ09; Beloborodov 2008). Although the bulk of the disc is initially driven neutron rich by degenerate electron captures, Fig. 2 shows that a large portion of the mass remaining at late times has $Y_e \gtrsim 0.5$. As discussed in Section 1.1, the disc composition is driven proton rich by the powerful $v_e$ flux from the central proto-NS.

We have performed similar calculations for other disc parameters in order to study the sensitivity of our results to the disc’s initial properties (see Table 1). For example, the discs found by D06 and D07 are fairly massive, but a less massive disc would result from a slower rotating WD. Thus, we have also performed calculations using the same initial density profile from D06, but normalizing the total initial disc mass to $M_{\nu,n} = 0.01$ M$_\odot$ instead of 0.1 M$_\odot$. We have also calculated models with $\Sigma(r, t = 0), L_e(t)$ and $\langle \epsilon_e \rangle(t)$ taken from the rapidly rotating, 1.92 M$_\odot$ WD model of D06, which formed a more massive initial disc ($M_{\nu,n} \approx 0.5$ M$_\odot$). In the high $M_{\nu,n}$ case, we find that a significant fraction of the disc has exactly
Table 1. Late-time mass and composition of AIC discs.

| $M_{d,0}$ (M⊙) | Model | α | $M_{ad}$ (M⊙) | $M_{ni}$ (M⊙) |
|----------------|-------|---|---------------|---------------|
| 0.1            | D06(1.46 M⊙ WD) | 0.3 | 0.030 | 0.021 |
| 0.1            | D06(1.46 M⊙ WD) | 0.03 | 0.021 | 8.9 × 10^{-3} |
| 0.01           | D06(1.46 M⊙ WD) | 0.3 | 4.3 × 10^{-3} | 1.8 × 10^{-3} |
| 0.5            | D06(1.92 M⊙ WD) | 0.3 | 0.13 | 0.066 |
| 0.5            | D06(1.92 M⊙ WD) | 0.03 | 0.11 | 0.051 |

(a) Initial disc mass, (b) model used for the initial surface density $\Sigma(r, t = 0)$ and the proto-NS neutrino luminosities and energies, (c) mass remaining when the disc is 90 per cent advective, (d) mass with $0.485 < Y_c < 0.6$ when the disc is 90 per cent advective.

$Y_c = 0.5$ at freeze out. This occurs because in this model $\alpha$-particle formation precedes weak freeze out; thus, any excess free neutrons or protons are almost immediately captured into $\alpha$-particles (which have $Y_c = 0.5$) following a neutrino absorption (the ‘alpha-effect’) of Fuller & Meyer (1995).

In Table 1, we summarize our results for the total mass ($M_{ad}$) and the total mass with $0.485 \lesssim Y_c \lesssim 0.6$ ($M_{ni}$) when the disc has become 90 per cent advective. The latter is labelled $M_{ni}$ because a significant fraction of the mass ejected with $Y_c$ in this range produces $^{56}\text{Ni}$ (see Section 3). In all cases, we find that a significant fraction ($\gtrsim 50$ per cent) of the mass remaining at late times has $Y_c \gtrsim 0.5$. Thus, under a variety of conditions, AIC leads to ~10 per cent of the initial disc being ejected as moderately proton-rich material.

3 OPTICAL TRANSIENTS FROM AIC

In some ways, AIC can be considered a ‘failed’ Type Ia SN because it does not produce ~1 M⊙ of shock-heated ejecta. Indeed, the total $^{56}\text{Ni}$ mass ejected from the proto-NS is $\lesssim 10^{-3}$ M⊙ (e.g. Woosley & Baron 1992; D06; D07; Metzger, Thompson & Quataert 2008b), rendering isolated NS birth practically invisible. In the presence of rotation, however, this failure is not complete. In Section 2, we showed that if an accretion disc forms around the proto-NS during collapse, a natural consequence of the disc’s subsequent evolution is that up to $\sim 10^{-2}$ M⊙ of material with $Y_c \gtrsim 0.5$ is ejected within a few seconds after core bounce.

When hyperaccreting discs become radiatively inefficient and begin forming $\alpha$-particles, gas and radiation pressure in the mid-plane are comparable and the entropy is $S \sim 10^{13} \text{baryon}^{-1}$. Subsequent outflows from the disc are likely to possess a comparable entropy because heating that occurs as matter is unbound (e.g. due to $\alpha$-particle formation) is unlikely to deposit more than a few $k_B$ baryon^{-1}. A typical outflow speed is $v \sim 0.1–0.2 \text{c}$: this is approximately the escape speed from the disc when it becomes advective and a similar speed is attained from the $\sim 9 \text{MeV} \text{ baryon}^{-1}$ released in producing Fe-peak elements.

Matter with $0.485 \lesssim Y_c \lesssim 0.6$ that begins inNSE at low entropy and expands adiabatically at mildly relativistic speeds is primarily synthesized into $^{56}\text{Ni}$ (e.g. Pruet, Thompson & Hoffman 2004) because $^{56}\text{Ni}$ is favoured in NSE under proton-rich conditions (Seitenzahl et al. 2008). Our results in Fig. 2 and Table 1 therefore imply that AIC outflows yield a typical Ni mass $M_{ni} \sim 10^{-2}$ M⊙ if the rotating 1.46 M⊙ WD progenitors used in the calculations of D06 and D07 are representative. The decay of this Ni can reheat the (adiabatically cooled) ejecta sufficiently to produce detectable transient emission once the outflow expands sufficiently that photons can diffuse out (Arnett 1982).

4 DETECTION PROSPECTS

Because AIC has never been observationally identified, its rate is uncertain. Binary population synthesis models predict galactic rates in the range $R_{\text{AIC}} \sim 10^{-4}–10^{-3}$ yr\(^{-1}\) (e.g. Yungelson & Livio 1998). The amount of highly neutron-rich material ejected per AIC can also be used to constrain the AIC rate (Hartmann, Woosley & El Eid 1985; Fryer et al. 1999): if the $\sim 3 \times 10^{-3}$ M⊙ of ejected material with $Y_c \lesssim 0.4$ in the calculations of D06 is representative, $R_{\text{AIC}}$ must be $\lesssim 10^{-4}$ yr\(^{-1}\) so as not to overproduce the abundances of neutron-rich isotopes in our Solar system. If the proto-NS forms with a strong large-scale magnetic field $\sim 10^{15}$ G, then a much larger quantity $\sim 0.1$ M⊙ of neutron-rich material is ejected by the proto-NS’s magneto-centrifugally driven wind (D07; Metzger et al. 2008b) and the AIC rate is constrained to be even lower.

Assuming that the AIC rate is proportional to the blue stellar luminosity (Phinney 1991), a Galactic rate of $R_{\text{AIC}} \equiv 10^{-4}–10^{-3}$ yr\(^{-1}\)
corresponds to a volumetric rate of $10^{-4} R_{\odot}$ Mpc$^{-1}$ yr$^{-1}$. For the 1.46 M$_{\odot}$ WD model of D06, we predict that AIC produces an optical transient with a peak luminosity $\sim 10^{41}$ erg s$^{-1}$ (Fig. 3). This corresponds to a maximum detection distance of 650(410)[100] Mpc for a limiting magnitude of 25(24)[21]. Thus, the PanSTARRs Medium Deep Survey (MDS), which covers $\sim 84$ deg$^2$ in $g$ and $r$ down to AB magnitude 25, will detect $\sim 2.4 R_{\odot}$ yr$^{-1}$. The Palomar Transient Factory (PTF), which surveys $\sim 8000$ deg$^2$ to a limiting AB magnitude of 21, will detect $0.9 R_{\odot}$ yr$^{-1}$. Thus, if $R_{\text{AIC}} \sim 10^{-4}$ yr$^{-1}$, MDS and PTF should detect a few AIC events per year. Prospects for detection are much better with the Large Synoptic Survey Telescope, which will image the entire sky down to a limiting magnitude $\sim 24.5$ every 3–4 nights and should detect AIC events at a rate $\sim 800 R_{\odot}$ yr$^{-1}$. The short, few-day-time-scale of the transient requires observations with high cadence. Shallow, wider-area surveys such as PTF might be sufficient because separating the transient from its host galaxy’s light may be difficult for high-redshift events.

The detectability of an optical transient is sensitive to whether there is sufficient mass in a remnant disc, which in turn depends on the angular momentum distribution of the WD progenitor. The very massive AIC progenitor from D06 (1.92 M$_{\odot}$) is consistent with the large shear found by Yoon & Langer (2004, 2005), who studied angular momentum transport by Kelvin–Helmholtz instabilities during accretion. Piro (2008) argued, however, that large magnetic stresses impose nearly solid-body rotation, which prevents the WD from reach a mass much over the (non-rotating) Chandrasekhar limit (this may also be the case for purely hydrodynamic instabilities; see Saio & Nomoto 2004). More work is needed to understand the possible progenitors of AIC, and the resulting implications for the accretion disc formed during collapse, taking into account the effects of magnetic stresses and compositional gradients on the redistribution of angular momentum during accretion.

Other thermal transients are predicted to occur in nature on $\sim$day time-scales, such as ‘Ia’ SNe due to unstable thermonuclear He flashes from WD binaries (Bildsten et al. 2007). Although such events may originate from a similar stellar population and could be confused with AIC, a defining characteristic of AIC outflows is that most of the ejected material is processed through low-entropy NSE. As a result, we expect few intermediate-mass elements to be produced (such as O, Ca and Mg), although small amounts of He may be present. Thus, the Ni-rich explosions produced by AIC should be spectroscopically distinct from other types of SN identified to date.

The time-scale, velocity and composition of the AIC transient can be modified if significant amounts of WD material remain at large radii following AIC, either because the WD is still collapsing at late times or because some matter remains centrifugally supported in a remnant disc created by a super-M$_{\odot}$ WD–WD merger (e.g. Yoon et al. 2007). In the later case, depending on the mass and composition of the WD binary, up to $\sim 0.1$ M$_{\odot}$ in C, O, Ne or He could remain at a radius $R \sim 10^8$ cm. Once the Ni wind (with energy $\sim 10^{50}$ erg) impacts this remnant material, it will (1) shock heat the material to a few times $10^4$ K, which may synthesize intermediate-mass elements (Woosley, Heger & Weaver 2002), but will not disassociate the Ni and will likely leave some unburned WD material and (2) blow the ejecta, from $\sim 0.1$ c to a few thousand km s$^{-1}$. A combination of slower ejecta and higher opacity would also lengthen the duration of the light curve by a factor of $\sim 10$ from that shown in Fig. 3, making such an event easier to detect. Indeed, such an ‘enshrouded AIC’ model represents a possible explanation for sub-luminous, sub-Chandrasekhar Type I SNe, such as 2008ha, which rose to its peak brightness in only $\sim 10$ d.

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