1. INTRODUCTION

The *Fermi* Gamma-ray Space Telescope was launched on 2008 June 11 on a mission to study the universe at high energies. One of *Fermi*’s highest priority objectives is to help explain the physical mechanisms responsible for the powerful high-energy emission from gamma-ray bursts (GRBs). The *Fermi* Gamma-ray Burst Monitor (GBM) supports that goal by detecting and measuring the prompt emission from GRBs and providing quick notification to *Fermi*’s main instrument, the Large Area Telescope (LAT), as well as to ground-based observers. The broad energy coverage of the GBM (8 keV–40 MeV) and the LAT (30 MeV–300 GeV) combine to measure the emission spectra of sufficiently bright GRBs over more than seven decades of energy.

The onboard GBM trigger system for detecting GRBs was first enabled on 2008 July 12. In this paper, we provide a catalog of GRBs that triggered the GBM during its first two years of operation. During this time the instrument burst detection algorithms were triggered 908 times: 492 of these are classified as triggered GRBs: sky location, duration, peak flux, and fluence. This catalog is an official product of the *Fermi* GBM science team, and the data files containing the complete results are available from the High-Energy Astrophysics Science Archive Research Center.

*Key words:* catalogs – gamma-ray burst: general

*Online-only material:* machine-readable tables

GBM comprises 12 NaI scintillation detectors and 2 BGO scintillation detectors. The NaI detectors are 0.5 inch thick, 5 inch diameter and operate in the energy range 8 keV to 1 MeV. The performance of the NaI detectors at low energies is significantly enhanced by the use of beryllium entrance windows. The BGO detectors are 5 inch thick by 5 inch diameter and operate in the energy range 8 keV to 80 MeV.

Detailed studies of various GBM GRB subsamples have been presented elsewhere (Guiriec et al. 2010; Ghirlanda et al. 2010, 2011; Zhang et al. 2011). Spectral characteristics derived from a standard analysis are described in a companion catalog (Goldstein et al. 2012). Observations using non-standard data types and techniques developed post-launch are also described by A. Abdo et al. (2012, in preparation).

The GBM design is largely based on the Burst and Transient Source Experiment (BATSE) on the *Compton Gamma Ray Observatory (CGRO)* (Fishman et al. 1989), which operated from 1991 to 2000. Both instruments employ multiple sodium iodide (NaI) detectors to achieve full sky field of view, have on board burst triggering capability, and use relative count rates to obtain approximate directions to bursts. GBM also includes two bismuth germanate (BGO) detectors that can better detect higher energy photons. BATSE, with significantly larger NaI detectors, had better sensitivity, while GBM has a broader energy range and higher data rate.

This catalog summarizes some basic characteristics of the triggered GRBs: sky location, duration, peak flux, and fluence. Spectral characteristics derived from a standard analysis are described in a companion catalog (Goldstein et al. 2012). Detailed studies of various GBM GRB subsamples have been presented elsewhere (Guiriec et al. 2010; Ghirlanda et al. 2010, 2011; Ly et al. 2010; Bissaldi et al. 2011; Gruber et al. 2011a; Nava et al. 2011a, 2011b; Zhang et al. 2011).

2. INSTRUMENTATION

GBM comprises 12 NaI scintillation detectors and 2 BGO scintillation detectors. The NaI detectors are 0.5 inch thick by 5 inch diameter and operate in the energy range 8 keV to 1 MeV. The performance of the NaI detectors at low energies is significantly enhanced by the use of beryllium entrance windows. Their positions and orientations on the spacecraft permit burst localization over the entire sky (unocculted by the Earth). The BGO detectors are 5 inch thick by 5 inch diameter and operate in the 200 keV to 40 MeV energy range. They are located on opposite sides of the spacecraft so that at least one of them is illuminated from any direction. The GBM flight software (FSW) continuously monitors the detector...
count rates to detect GRBs and other short-timescale transients, computes their location on the sky, provides a preliminary classification, and promptly notifies the LAT of their occurrence. Once triggered, the FSW initiates prompt transmission of a subset of the data to the ground for quick-look analysis and notification of ground-based instruments via the Gamma-ray Coordinates Network (GCN). The instrument is described in more detail by Meegan et al. (2009).

The GBM trigger algorithms operate on background-subtracted count rates over a programmable range of timescales (from a minimum of 16 ms to a maximum of 16,384 s; currently the longest is 4 s) and in four different energy ranges (currently 25–50 keV, 50–300 keV, >100 keV, and >300 keV). This is the primary method for detection of GRBs, TGFs, SGRs, and solar flares.

*Fermi’s* use of the *Ku* band for downlink of science data allows GBM to have a higher data rate than BATSE, which generally translates to better time and energy resolution. Outside of a trigger, the GBM continuously transmits two types of science data: continuous time (CTIME) and continuous spectroscopy (CSPEC). The CTIME data have finer time resolution (nominally 256 ms but configurable from 64 ms to 1,024 s in units of 64 ms) and coarse energy resolution (eight channels). The CSPEC data have the full energy resolution (128 channels) but coarser time resolution (nominally 4,096 s but configurable from 1,024 s to 32,768 s in units of 1,024 s). In this mode time-tagged event (TTE) data are also produced but not transmitted to the ground. The TTE data consist of individual detector events, each tagged with arrival time (2 μs resolution, accurate to ∼10 μs), energy (128 channels), and detector number. These non-triggered TTE data are temporarily stored on board in a ring buffer with a capacity of 512,000 events, which lasts for 25–30 s at typical background rates.

Upon entering trigger mode, the FSW speeds up CTIME resolution to 64 ms and CSPEC resolution to 1,024 s. In addition, the TTE data are transmitted directly to the science data bus instead of being stored in the ring buffer (the contents of which are frozen for later downlink). The production of prompt TTE lists for approximately 300 s, at which point the contents of the pre-trigger ring buffer are dumped to the science data bus. After an additional 300 s, the CTIME and CSPEC data are returned to their non-triggered time resolution and the FSW exits trigger mode.

The basic trigger design follows that used for BATSE: to trigger, two or more detectors must have a statistically significant rate increase above the background rate. Requiring two detectors to be above their respective thresholds makes triggering on statistical fluctuations less likely, and much more importantly, it suppresses triggering due to non-astrophysical effects that appear in only one detector, such as phosphorescence spikes. Each algorithm has its own threshold setting, configurable from 0.1σ to 25.5σ in units of 0.1σ. The background model is a trailing average of the detector data. Whereas BATSE used three trigger algorithms (a single energy range, usually 50–300 keV, and three timescales, 64 ms, 256 ms, and 1,024 s), the GBM FSW supports up to 119 trigger algorithms. A given algorithm is defined by its timescale, offset, and energy range. The offset is a value in milliseconds by which the timescale is shifted. Running overlapping accumulations for a given combination of timescale and energy range provides some improvement in trigger sensitivity (Band 2002; Band et al. 2004).

While in trigger mode, the FSW continues to monitor other enabled algorithms and records all instances where a given algorithm would also have triggered on the same burst. A special *compute* mode is available in which an algorithm is monitored during a trigger (initiated by any other algorithm) to determine if its threshold is exceeded but no triggers can actually be initiated by that algorithm.

The FSW classifies triggers based on a number of criteria, including the event localization, spectral hardness, and the spacecraft geomagnetic latitude (McIlwain L coordinate). A Bayesian approach is used to assign identification probabilities for various event classes, including GRB, solar flare, SGR, particle precipitation, and known transient sources. Classification of TGFs is a special case: early in the mission it was observed that these events trigger NaI detectors exclusively on the combination of shortest timescale (0.016 s) and one of the high-energy ranges (>100 keV or >300 keV).

While in trigger mode, the FSW continues to monitor the detector rates on various timescales and, based on improved statistical significance, updates of the localization and classification may be generated and transmitted in the quick-look science data.

3. IN-ORBIT OPERATIONS

3.1. Trigger Criteria

GBM triggering has been enabled continuously since 2008 July 12, except during South Atlantic Anomaly (SAA) passages and a few brief intervals when FSW upgrades were being installed. Early in the mission the trigger algorithms used only data from the NaI detectors. This turned out to be a limitation for TGFs because they are spectrally much harder than GRBs. In 2009 November the FSW was revised to add trigger algorithms that use the BGO detectors, which significantly improved the GBM sensitivity for detecting TGFs.

In orbit, the GBM has enabled 71 trigger algorithms, 5 of which are TGF-specific algorithms that use the BGO detectors alone or combined with NaI detectors. No GRB has ever triggered only on a TGF-specific algorithm, so they are not discussed further in this paper. Table 1 summarizes the 66 algorithms relevant for this catalog, the times during which they were enabled, and the history of their threshold settings.

During the first year most of the algorithms were enabled and a few minor adjustments made to their thresholds. Exceptions are the 16 s algorithms (numbers 20, 21, 41, and 42), which were deemed too sensitive to background variations and disabled after ~2 weeks. After nearly a year’s experience, it was judged that a large number of the algorithms were of dubious value because they never triggered on an event that did not also trigger another algorithm. This list included all of the 8 s algorithms as well as most of the algorithms not operating in the standard BATSE energy range (50–300 keV). In order to ease the computational burden on the FSW, these algorithms were disabled in early July of 2009. The configuration in the rightmost column of Table 1 remained the same throughout the remainder of the period covered by this catalog.

Direct comparisons between the BATSE and GBM data sets are possible since GBM includes the same three trigger algorithms used by BATSE (64, 256, and 1024 ms timescales in the 50–300 keV energy range). Depending on which set of the half-bin offset algorithms are considered as the BATSE algorithms, we find that either 405 or 408 of the 491 GRBs would have triggered GBM, in agreement with pre-launch
Table 1
Trigger Criteria History

| Algorithm Number | Timescale (ms) | Offset (ms) | Channels | Energy (keV) | Energy Threshold (0.1σ) |
|------------------|----------------|-------------|----------|-------------|-------------------------|
| 1                | 16             | 0           | 3–4      | 50–300      | 75                      |
| 2                | 32             | 0           | 3–4      | 50–300      | 75                      |
| 3                | 32             | 16          | 3–4      | 50–300      | 75                      |
| 4                | 64             | 0           | 3–4      | 50–300      | 45                      |
| 5                | 64             | 32          | 3–4      | 50–300      | 45                      |
| 6                | 128            | 0           | 3–4      | 50–300      | 45                      |
| 7                | 128            | 64          | 3–4      | 50–300      | 45                      |
| 8                | 256            | 0           | 3–4      | 50–300      | 45                      |
| 9                | 256            | 128         | 3–4      | 50–300      | 45                      |
| 10               | 512            | 0           | 3–4      | 50–300      | 45                      |
| 11               | 512            | 256         | 3–4      | 50–300      | 45                      |
| 12               | 1024           | 0           | 3–4      | 50–300      | 45                      |
| 13               | 1024           | 512         | 3–4      | 50–300      | 45                      |
| 14               | 2048           | 0           | 3–4      | 50–300      | 45                      |
| 15               | 2048           | 1024        | 3–4      | 50–300      | 45                      |
| 16               | 4096           | 0           | 3–4      | 50–300      | 45                      |
| 17               | 4096           | 2048        | 3–4      | 50–300      | 45                      |
| 18               | 8192           | 0           | 3–4      | 50–300      | C 50                    |
| 19               | 8192           | 4096        | 3–4      | 50–300      | C 50                    |
| 20               | 16384          | 0           | 3–4      | 50–300      | C 50 D                  |
| 21               | 16384          | 8192        | 3–4      | 50–300      | C 50 D                  |
| 22               | 16             | 0           | 2–2      | 25–50       | D 80                    |
| 23               | 32             | 0           | 2–2      | 25–50       | D 80                    |
| 24               | 32             | 16          | 2–2      | 25–50       | D 80                    |
| 25               | 64             | 0           | 2–2      | 25–50       | D 55                    |
| 26               | 64             | 32          | 2–2      | 25–50       | D 55                    |
| 27               | 128            | 0           | 2–2      | 25–50       | D 55                    |
| 28               | 128            | 64          | 2–2      | 25–50       | D 55                    |
| 29               | 256            | 0           | 2–2      | 25–50       | D 55                    |
| 30               | 256            | 128         | 2–2      | 25–50       | D 55                    |
| 31               | 512            | 0           | 2–2      | 25–50       | D 55                    |
| 32               | 512            | 256         | 2–2      | 25–50       | D 55                    |
| 33               | 1024           | 0           | 2–2      | 25–50       | D 55                    |
| 34               | 1024           | 512         | 2–2      | 25–50       | D 55                    |
| 35               | 2048           | 0           | 2–2      | 25–50       | D 55                    |
| 36               | 2048           | 1024        | 2–2      | 25–50       | D 55                    |
| 37               | 4096           | 0           | 2–2      | 25–50       | D 55                    |
| 38               | 4096           | 2048        | 2–2      | 25–50       | D 55                    |
| 39               | 8192           | 0           | 2–2      | 25–50       | D 55                    |
| 40               | 8192           | 4096        | 2–2      | 25–50       | D 55                    |
| 41               | 16384          | 0           | 2–2      | 25–50       | D 65 D                   |
| 42               | 16384          | 8192        | 2–2      | 25–50       | D 65 D                   |
| 43               | 16             | 0           | 5–7      | >300        | D 80                    |
| 44               | 32             | 0           | 5–7      | >300        | D 80                    |
| 45               | 32             | 16          | 5–7      | >300        | D 80                    |
| 46               | 64             | 0           | 5–7      | >300        | D 55                    |
| 47               | 64             | 32          | 5–7      | >300        | D 55                    |
| 48               | 128            | 0           | 5–7      | >300        | D 55                    |
| 49               | 128            | 64          | 5–7      | >300        | D 55                    |
| 50               | 16             | 0           | 4–7      | >100        | D 80                    |
| 51               | 32             | 0           | 4–7      | >100        | D 80                    |
| 52               | 32             | 16          | 4–7      | >100        | D 80                    |
| 53               | 64             | 0           | 4–7      | >100        | D 55                    |
| 54               | 64             | 32          | 4–7      | >100        | D 55                    |
| 55               | 128            | 0           | 4–7      | >100        | D 55                    |
| 56               | 128            | 64          | 4–7      | >100        | D 55                    |
| 57               | 256            | 0           | 4–7      | >100        | D 55                    |
| 58               | 256            | 128         | 4–7      | >100        | D 55                    |
| 59               | 512            | 0           | 4–7      | >100        | D 55                    |
| 60               | 512            | 256         | 4–7      | >100        | D 55                    |
| 61               | 1024           | 0           | 4–7      | >100        | D 55                    |
| 62               | 1024           | 512         | 4–7      | >100        | D 55                    |
| 63               | 2048           | 0           | 4–7      | >100        | D 55                    |
| 64               | 2048           | 1024        | 4–7      | >100        | D 55                    |
estimates of 200 bursts/year (Meegan et al. 2007). Using both sets of overlapping windows raises the total number of triggers to 423, leaving 68 events that did not trigger on any BATSE-style algorithms. Of these, 63 (93%) triggered the longer (>1.024 s) timescale algorithms in the 47–291 keV energy range, 3 triggered only the 512 ms algorithm in 47–291 keV, and 2 triggered only in the 23–47 keV energy range (1 each in 64 ms and 2.048 s). Thus, the apparent improvement in trigger sensitivity relative to BATSE is attributable mainly to GBM’s additional longer trigger timescales.

3.2. Quick-look Analysis

As described above, quick-look data are generated during trigger mode and promptly transmitted to the ground. For events classified as GRBs, the FSW-generated localization and classification information is further distributed via GCN notices. Also for GRBs, additional GCN notices containing ground-generated localizations are produced and distributed automatically. The GBM location algorithm is an adaptation of the method developed for BATSE (Pendleton et al. 1999).

Table 1

| Algorithm | Timescale | Offset | Channels | Energy | 2008 Jul 11 | 2008 Jul 14 | 2008 Aug 1 | 2009 May 8 | 2009 Jul 2 |
|-----------|-----------|--------|----------|--------|-------------|-------------|------------|------------|------------|
| Number    | (ms)      | (ms)   | (keV)    |        |             |             |            |            |            |
| 65        | 4096      | 0      | 4–7      | >100   | D           |             |            |            | D          |
| 66        | 4096      | 2048   | 4–7      | >100   | D           |             |            |            | D          |

Note. a Symbol “::” indicates no change from previous setting; “C” indicates that the algorithm is in compute mode (see the text); and “D” indicates that the algorithm is disabled.

Table 2

| Trigger ID | GRB Name | Time (UT) | α (°) | δ (°) | Error (°) | Location | Algorithm | Timescale (ms) | Energy (keV) | Other Detections |
|------------|----------|-----------|-------|-------|-----------|----------|-----------|----------------|--------------|-----------------|
| bn080714B  | GRB 080714B | 02:04:12.0534 | 41.9  | 8.5   | 7.5       | Fermi-GBM | 10        | 512           | 47–291       | IA              |
| bn080714C  | GRB 080714C | 10:12:01.8376 | 187.5 | −74   | 8.7       | Fermi-GBM | 17        | 4096          | 47–291       | IA              |
| bn080717A  | GRB 080717A | 17:52:54.0234 | 188.1 | −60.2 | 0         | Swift     | 13        | 1024          | 47–291       | S, K, IA, R     |
| bn080719A  | GRB 080719A | 12:41:40.9578 | 153.2 | −61.3 | 13.8      | Fermi-GBM | 16        | 4096          | 47–291       | IA              |
| bn080720A  | GRB 080720A | 07:35:35.5476 | 98.5  | −43.9 | 4.8       | Fermi-GBM | 19        | 8192          | 47–291       | IA              |
| bn080723B  | GRB 080723B | 13:22:21.3751 | 176.8 | −60.2 | 0         | Swift     | 8         | 256           | 47–291       | IB, IA, K, SA, AM |
| bn080723C  | GRB 080723C | 21:55:23.0583 | 113.3 | −19.7 | 9.9       | Fermi-GBM | 5         | 64            | 47–291       | W               |
| bn080723D  | GRB 080723D | 23:37:42.7083 | 105.3 | 71.1  | 1         | Fermi-GBM | 11        | 512           | 47–291       | K, W*, IA        |

Notes.

a Bursts with Trigger ID and GRB Name in italics have significant emission in at least one BGO detector (see the text).
b Other instrument detections: K: Konus-Wind, K-RF: Konus-RF, S: Swift, IA: INTEGRAL SPI-ACS, IB: INTEGRAL Burst Alert System, W: Suzaku-WAM, R: RHESSI, M: MAXI, SA: SuperAGILE, AM: AGILE-MCAL, A: AGILE, L: Fermi LAT
c GRB 091024A triggered GBM twice.

Figure 1. Sky distribution of GBM triggered GRBs in celestial coordinates. Crosses indicate long GRBs ($T_{90} > 2$ s); asterisks indicate short GRBs.

Fermi GBM GRBs in first two years of operation

491 GBM GRBs
4 400 Long
x 88 Short
o 3 No measured duration
### Table 3
**GRB Durations (50–300 keV)**

| Trigger ID | Detectors Used | $T_{50}$ | $T_{90}$ Start | $T_{50}$ Start |
|------------|----------------|---------|----------------|---------------|
| bn080714086 | 3+4+8         | 5.376 ± 2.360 | −0.768 | 2.816 ± 0.810 | −0.256 |
| bn080714425 | 0+9+10       | 40.192 ± 1.145 | −4.352 | 11.776 ± 1.619 | −1.280 |
| bn080714745 | 5             | 59.649 ± 11.276 | −0.512 | 25.088 ± 7.940 | 2.560 |
| bn080715950 | 0+1+2+9+10   | 7.872 ± 0.272 | 0.128 | 6.144 ± 0.264 | 1.088 |
| bn080717543 | 2+10          | 36.609 ± 2.985 | −5.376 | 13.056 ± 0.810 | 1.024 |
| bn080719529 | 6+7+9         | 16.128 ± 17.887 | −4.352 | 8.448 ± 1.280 | −2.048 |
| bn080720316a | ...          | ...       | ...     | ...           | ...   |
| bn080723557 | 4             | 58.369 ± 1.985 | 2.368 | 40.513 ± 0.231 | 14.208 |
| bn080723913 | 0+1+3        | 0.192 ± 0.345 | −0.064 | 0.064 ± 0.143 | −0.064 |
| bn080723985 | 2+5           | 42.817 ± 0.659 | 3.072 | 25.280 ± 0.405 | 12.160 |

#### Notes.
- a Data problems precluded duration analysis.
- b Used TTE binned at 32 ms.
- c Partial Earth occultation is likely; durations are lower limits.
- d Possible precursor at $\sim T_0 - 120$ s.
- e Data cutoff while burst in progress; durations are lower limits.
- f SAA entry at $T_0 + 83$ s; durations are lower limits.
- g GRB 091024 triggered GBM twice.
- h Too weak to measure durations; visual duration is $\sim 0.025$ s.
- i Possible contamination due to emergence of Crab & A0535+26 from Earth occultation.
- j Used TTE binned at 16 ms.

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

### Table 4
**GRB Fluence and Peak Flux (10–1000 keV)**

| Trigger ID | Fluence (erg cm$^{-2}$) | PF64 (photons cm$^{-2}$ s$^{-1}$) | PF256 (photons cm$^{-2}$ s$^{-1}$) | PF1024 (photons cm$^{-2}$ s$^{-1}$) |
|------------|-------------------------|----------------------------------|----------------------------------|------------------------------------|
| bn080714086 | 6.76E-07 ± 4.10E-08     | 3.82 ± 1.06                     | 2.24 ± 0.36                     | 1.54 ± 0.18                        |
| bn080714425 | 1.81E-06 ± 2.10E-08     | 4.00 ± 1.45                     | 2.96 ± 0.46                     | 2.02 ± 0.21                        |
| bn080714745 | 6.33E-06 ± 1.41E-07     | 8.89 ± 1.61                     | 7.78 ± 0.83                     | 6.93 ± 0.39                        |
| bn080715950 | 5.04E-06 ± 8.00E-08     | 19.42 ± 0.95                    | 13.58 ± 0.45                    | 9.91 ± 0.22                        |
| bn080717543 | 4.46E-06 ± 7.70E-08     | 6.24 ± 1.08                     | 3.43 ± 0.49                     | 2.89 ± 0.23                        |
| bn080719529 | 7.75E-07 ± 2.90E-08     | 2.77 ± 0.83                     | 1.77 ± 0.29                     | 1.12 ± 0.16                        |
| bn080720316a | ...                   | ...                             | ...                             | ...                                |
| bn080723557 | 7.22E-05 ± 2.54E-07     | 40.97 ± 2.24                    | 38.24 ± 1.09                    | 30.45 ± 0.49                       |
| bn080723913 | 1.34E-07 ± 1.40E-08     | 5.26 ± 0.70                     | 4.13 ± 0.32                     | 1.41 ± 0.13                        |
| bn080723985 | 3.08E-05 ± 2.07E-07     | 13.45 ± 1.24                    | 11.36 ± 0.60                    | 10.12 ± 0.28                       |

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

### Table 5
**GRB Fluence and Peak Flux (50–300 keV)**

| Trigger ID | Fluence (erg cm$^{-2}$) | PF64 (photons cm$^{-2}$ s$^{-1}$) | PF256 (photons cm$^{-2}$ s$^{-1}$) | PF1024 (photons cm$^{-2}$ s$^{-1}$) |
|------------|-------------------------|----------------------------------|----------------------------------|------------------------------------|
| bn080714086 | 3.54E-07 ± 1.70E-08     | 1.52 ± 0.74                     | 0.91 ± 0.36                     | 0.43 ± 0.18                        |
| bn080714425 | 9.79E-07 ± 1.40E-08     | 1.03 ± 0.45                     | 0.71 ± 0.19                     | 0.46 ± 0.08                        |
| bn080714745 | 3.26E-06 ± 6.00E-08     | 4.41 ± 1.66                     | 3.27 ± 0.71                     | 2.82 ± 0.36                        |
| bn080715950 | 2.54E-06 ± 3.50E-08     | 10.70 ± 0.95                    | 6.61 ± 0.45                     | 3.83 ± 0.22                        |
| bn080717543 | 2.37E-06 ± 4.30E-08     | 2.14 ± 1.03                     | 1.30 ± 0.47                     | 1.05 ± 0.23                        |
| bn080719529 | 3.88E-07 ± 1.50E-08     | 0.59 ± 0.18                     | 0.32 ± 0.08                     | 0.23 ± 0.04                        |
| bn080720316a | ...                   | ...                             | ...                             | ...                                |
| bn080723557 | 3.92E-05 ± 1.15E-07     | 21.19 ± 1.79                    | 19.81 ± 1.09                    | 15.14 ± 0.48                       |
| bn080723913 | 7.50E-08 ± 5.00E-09     | 2.62 ± 0.66                     | 2.14 ± 0.32                     | 0.69 ± 0.13                        |
| bn080723985 | 1.57E-05 ± 1.07E-07     | 5.92 ± 1.23                     | 5.17 ± 0.54                     | 4.85 ± 0.28                        |

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Figure 2. Distribution of GRB durations in the 50–300 keV energy range. The upper plot shows $T_{50}$ and the lower plot shows $T_{90}$.

 FSW and ground locations use the count rates in all 12 NaI detectors to point back to a preferred direction on the sky by comparing observed rates to model rates and minimizing $\chi^2$. The model rates are a combination of counts that come directly from the source into the detector, counts from the source scattered in the spacecraft into the detectors, and counts from source photons that hit Earth’s atmosphere and are scattered into the detectors. All three of these components are a function of the source intensity, its spectrum, and the source-spacecraft geometry, with the final component also depending on the source-spacecraft-Earth geometry. For automated locations on board and on the ground, the background count rate subtracted from the observed counts is an average over a 16 s interval before the burst trigger occurred. However, the ground automated localizations differ from the flight locations in several ways:

1. Although the two decision making processes use the same rate data type, they run independently with different criteria and do not necessarily use rate data from the same time intervals.

2. The ground process has access to location tables generated with finer sky resolution (1°, compared to 5° for the FSW).

3. The ground process includes a more accurate treatment of atmospheric scattering (based on the actual orientation of the spacecraft with respect to Earth, whereas the FSW assumes zenith pointing for all model rates).

4. The ground process incorporates the spectrum of the source into the calculation of the expected rates by choosing one of three location tables based on the hardness of the burst as determined by the ratio of counts $>50$ keV to counts $<50$ keV.

The GBM team assigns a burst advocate (BA) to inspect the real-time data promptly and perform additional analysis as appropriate. Normally, the BA will generate additional localizations and optionally distribute these via the GCN (circulars were used during the time period of this catalog but currently GCN notices are used). These “human-in-the-loop” localizations use source and background time intervals and model fits selected by the user based on the entire quick-look data set, which extends from 200 s pre-trigger to 450 s post-trigger. The BAs typically run the location code several times, using different selections of time interval and/or background models, and select a best location using statistical error and goodness-of-fit criteria. This is particularly useful in verifying that separate pulses are consistent with the same sky location. The FSW classification is reviewed by the BA, usually in consultation with other GBM team members, and may be corrected based on inspection of the GBM quick-look data and/or additional information such as detection by another instrument.

4. CATALOG ANALYSIS

4.1. Burst Localization and Instrument Response

Determination of the approximate burst sky location is important because the other results reported in this catalog and
the companion spectroscopy catalog (Goldstein et al. 2012) require instrument response functions that are dependent on the direction of the burst relative to the detectors and to the Earth. Most of the burst locations reported in this catalog are the result of the manual on-ground analysis, typically by the BA as described in Section 3.2. If, however, the burst was also localized with better precision by another instrument (e.g., Swift or the LAT) that location was used instead to derive the GBM instrument response for the subsequent catalog analysis. A total of 76 bursts have locations from Swift; 63 of these triggered the Swift Burst Alert Telescope (BAT) and 5 more were detected in ground analysis of BAT data. The remaining eight were located by the Swift X-Ray Telescope (XRT) following detections of prompt emission by other instruments.

The accuracy of the GBM burst localizations was checked by comparing the independently derived GBM locations with a sample of higher precision locations obtained by other instruments for the same GRBs. Using 127 bursts with known locations (some of which occurred after the end of the current catalog), we find that the true GBM human-in-the-loop location errors are best described by combining the statistical error in quadrature with a systematic error, where the current best model for systematic errors is $T_{90} = 2$ s as the distinguishing criterion. In each plot a power law with a slope of $-3/2$ (dashed line) is drawn to guide the eye.
radius, 93 (73%) are within twice the $1\sigma$ radius, and 107 (84%) are within three times the $1\sigma$ radius. A more detailed analysis of the GBM location errors is in progress and will be reported later (V. Connaughton et al. 2012, in preparation).

Figure 1 shows the sky distribution of the GBM-detected GRBs in celestial coordinates. The large-scale isotropic distribution is well known from BATSE observations (Briggs et al. 1996) and the GBM distribution appears to be consistent with this.

4.2. Duration, Peak Flux, and Fluence

In addition to the burst locations, the present paper reports various measures of the duration, peak flux, and fluence of each burst, with a few exceptions due to analysis difficulties such as incomplete data or background interference. The burst durations $T_{90}$ and $T_{90}^3$ were computed in the 50–300 keV energy range. The fluence for each burst was computed in two energy ranges: 50–300 keV and 10–1000 keV. Peak fluxes for each burst were computed in these same energy ranges and for three different timescales: 64 ms, 256 ms, and 1024 ms.

Burst durations were determined using a method similar to that developed for BATSE (Kouveliotou et al. 1993; Koshut et al. 1996). However, in the BATSE analysis all quantities were derived from the counts directly, whereas in the present analysis the counts spectrum in each time bin is deconvolved and the durations are computed from the time history of fitted photon spectra. Peak fluxes and fluences are naturally obtained in the same analysis, using the same choices of detector subset, source and background intervals, and background model fits. Unlike CGRO, which was inertially pointed in the same orientation for weeks at a time, the Fermi observatory has been operated in an all-sky survey mode during the period covered by this catalog. To optimize sky coverage, the spacecraft sweeps its $z$-axis across the sky at a specified angle perpendicular to the orbit plane (currently 50$^\circ$), rocking on alternate orbits above and below the orbit plane by the specified angle. Within a given orbit Fermi also executes a slow roll about the $z$-axis to maintain optimal orientation of the solar panels with respect to the Sun.

In the energy range of interest, the response of a given GBM NaI detector varies approximately as the cosine of the angle between the source direction and the detector axis. Therefore as Fermi slews, the detector to source angle changes and hence so does the response, with the rate of change being different for each detector. This was not a factor for BATSE, where inertial pointing kept the source to detector angles constant for extended periods. The most accurate correction for the response changing over time is a spectral deconvolution (assuming that the changing response is tracked correctly), so we have adopted the deconvolved flux history as our basis for determining the duration parameters. Furthermore, comparisons from burst to burst of the peak flux and fluence will not be compromised by differences in response arising from different source angles, as they would be for the raw counts. Finally, a fluence that is obtained by integrating a deconvolved flux history incorporates spectral evolution throughout the event in a way that summed counts can never attain, due to the loss of temporal information from the summing. The Appendix describes the procedure in some detail.

For each burst, a set of NaI detectors was chosen with good source viewing angles (<60$^\circ$) and no apparent blockage by any other element of the spacecraft. For the majority of bursts the GBM CTIME data, which have 256 ms time resolution pre-trigger and 64 ms resolution post-trigger, were used. TTE data were used for bursts where at least one of the peak fluxes occurs at or before the trigger time, which happens for many short bursts and a few longer ones. A limitation is that the pre-trigger TTE data typically span at most 30 s, which in some cases was not enough for computing the background and for some long bursts included significant burst emission. In such cases, the analysis was done with CTIME data. When using TTE data, which have 256 channels of energy resolution, it was often found that the deconvolution analysis is more robust if the 256 pulse-height channels were first summed into eight channels, as they would be for the raw counts. Finally, a fluence that is obtained by integrating a deconvolved flux history incorporates spectral evolution throughout the event in a way that summed counts can never attain, due to the loss of temporal information from the summing. The Appendix describes the procedure in some detail.

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5. CATALOG RESULTS

The catalog results can be accessed electronically through the High-Energy Astrophysics Science Archive Research Center (HEASARC) browse interface (http://heasarc.gsfc.nasa.gov/W3Browse/fermi/fermigbrst.html). Standard light curve plots for each burst can be viewed at http://gammaray.nsstc.nasa.gov/gbm/science/grbs/month_listings/. Here we provide tables that summarize selected parameters.

Table 2 lists the 492 triggers that were classified as GRBs. The GBM Trigger ID is shown along with a conventional
GRB name as defined by the GRB-observing community. For readers interested in the bursts with significant emission in the BGOs, the trigger ID and GRB name are highlighted in italics if emission in the BGO data (above 300 keV) is visible in the standard light curve plots.10 Note that the entire table is italics if emission in the BGO data (above 300 keV) is visible.

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6. DISCUSSION

Histograms of the \( T_{50} \) and \( T_{90} \) distributions are shown in Figure 2. Using the conventional division between the short and long GRB classes (\( T_{90} = 2 \) s), we find 88 (18%) of the 487 measured GRBs to be in the short classification. Within the quoted duration errors, the number of short GBM events ranges from 73 (15%) to 104 (21%). For comparison, the fraction of short events in the BATSE GRB catalog is 24%. The difference from BATSE is probably not due to a deficit of short events but rather to an excess of long events detected by the GBM’s longer timescale trigger algorithms (see Section 3.1).

The anti-correlation of spectral hardness with duration is well known from BATSE data (Kouveliotou et al. 1993) and a simple analysis shows that such an effect is also present in the GBM catalog. Time-resolved spectral fits for each GRB are a byproduct of the duration analysis and those photon model fit parameters were used to derive a measure of average spectral hardness. Figure 3 shows scatter plots of hardness derived in this way as a function of duration for the two duration measures. Although the effect of the 64 ms duration quantization is noticeable, the anti-correlation of hardness versus duration is visibly evident in the GBM data.

Integral distributions of the peak fluxes are shown in Figures 4–6 for the three different timescales and separately for short and long GRBs. For the long GRBs, deviation from the \(-3/2\) power law that would be expected if the GRBs were spatially homogeneous occurs well above the GBM threshold at a flux value of \(~10\) photons \( \text{s}^{-1} \text{cm}^{-2} \). This is consistent with earlier BATSE measurements (Paciesas et al. 1999), which have much better statistics. For the short events the GBM data appear consistent with a homogeneous spatial distribution down to peak flux values around 1 \( \text{photons} \text{s}^{-1} \text{cm}^{-2} \) (50–300 keV), below which instrument threshold effects become dominant. The integral fluence distributions for the two energy intervals are shown in Figure 7.

7. SUMMARY

The first GBM catalog includes 491 cosmic GRBs that triggered GBM between 2008 July 12 and 2010 July 11. Compared to BATSE, GBM has a higher threshold for burst detection but this is somewhat ameliorated by GBM’s additional range of trigger timescales (primarily the 4 s timescale) and, to a lesser extent, trigger energy ranges. The distribution of GBM durations is consistent with the well-known bimodality measured previously. The fraction of short GRBs in the GBM sample is somewhat smaller than detected by BATSE, which is attributed mainly to GBM’s ability to trigger on longer timescales.

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APPENDIX

COMPUTATION OF DURATION, PEAK FLUX, AND FLUENCE

For the catalog analysis, a standard calculation of durations, peak fluxes, and fluences was implemented as an add-on to the RMFIT software package that was developed for time-resolved analysis of BATSE GRB data but has been adapted for GBM and other instruments. For each burst, selections of the detectors and data types to be used for the catalog analysis were performed by the user based on uniform criteria as described in the main text. The remainder of this appendix describes the procedure and some caveats.
After the data have been selected for a given GRB and read into rmfit, a background model, polynomial in time, is computed separately for each detector, based on user-selected time intervals. The intervals normally include sections earlier and later than the evident burst emission, such that the background model can be interpolated through the entire time of burst activity. In some cases, the background selection includes quiescent portions between pulses of the GRB, if such regions are clearly identifiable. As a final preparation step, the interval for the duration analysis is chosen to cover the entire burst emission in time, as well as overlapping with the background selections. Figure 8 is an example CTIME light curve for GRB 081009A, showing the background and source selections used for the duration analysis.

The next step involves joint spectral fits using the selected detectors for each time bin in the selection, including the background regions before and after the burst. A poor estimate of the residual spectrum may result in very poor spectral fits, dominated by the default model parameters, so the selection of these values may be crucial, especially for weak bursts. In many cases, the duration estimate is more robust when the default parameters are set to values that are representative of the background itself (typical values are $E_{\text{peak}} = 70$ keV, $\alpha = -1.0$). Although this improves the duration analysis, it may introduce additional systematic errors in the computation of the fluences and peak fluxes.

The choice of photon model to fit is dictated by the sparse data statistics: the GBM TTE default time binning and CTIME post-trigger accumulation interval is 64 ms, resulting in average source counts that are an order of magnitude less than usually required for high-quality spectroscopy ($>45\sigma$). The model chosen for the catalog analysis is an exponentially cutoff power law, parameterized such that the characteristic energy ($E_{\text{peak}}$) is identical with the peak in $\nu F_\nu$ (the so-called COMP model in RMFIT). This model lacks a non-thermal high-energy power law, which is ideal, since it is precisely that parameter that would be least constrained by the sparse data at high energies. At the same time, it is desirable to constrain the three model parameters that describe the COMP photon model: amplitude, power-law index, and $E_{\text{peak}}$. Hence it is preferable to use data types that have few energy channels, so that there are better statistics in the channels at hand. CTIME, with eight energy channels, is usually the best choice, but the pre-trigger time resolution of 256 ms is a limitation for most short bursts as well as a few longer ones. In those cases TTE data may be preferred. Native TTE data have the full 128 channels available; however, for the duration analysis these are usually summed to match the energy channels of the corresponding CTIME data. 11 The data are fitted to the available CTIME energy channels, which cover the approximate energy range 10–1000 keV.

The goal in this analysis is to produce uncertainties in the flux determinations that are no worse than the statistical uncertainties due to the counts, while retaining the advantages of correcting for the detector response, which is only possible by doing a spectral fit. Just as important, we do not have to reproduce a detailed temporal spectral analysis for each burst, which would require summing the data over time until a significant sample has been accumulated and would also require better spectral resolution and more complicated spectral models. Instead, we only require that the spectral fit in each time bin be reasonably accurate over the energy bins used for the duration calculation (50–300 keV). Thus, it is unlikely that the presence of hard emission (such as a high-energy power law, as opposed to our choice of an exponential roll-off at high energies) in the data would make a significant contribution to the flux integral. The duration energy range is covered by the maximum in the detector response, so the spectral fit is best determined for those energies. The flux uncertainty is calculated from the covariance matrix

11 For the duration analysis, separate software is used to produce a new data type with eight energy channels, called CTTE, that is then read into rmfit. This is necessary because by design rmfit always uses the raw energy binning of the input data for spectral fitting.
Figure 9. Photon flux light curve for GRB 081009A (bn081009140) produced by the duration analysis. Data from NaI detectors 3, 4, 7, and 8 were used. Temporal resolution is the same as in the raw CTIME data. Vertical dotted lines are as described in the caption of Figure 10.

Figure 10. Duration plot for GRB 081009A (bn081009140) is an example of the analysis for a bright GRB. Data from NaI detectors 3, 4, 7, and 8 were used. Horizontal dotted lines are drawn at 5%, 25%, 75%, and 95% of the total fluence. Vertical dotted lines are drawn at the times corresponding to those same fluences, thereby defining the \( T_{50} \) and \( T_{90} \) intervals.

Figure 11. Duration plot for GRB 090531B (bn090531775) is an example of the analysis for a weak GRB. Data from NaI detectors 6, 7, and 9 were used. Dotted lines are as described in the caption for Figure 10.

The excess residuals are typically outside the 50–300 keV energy range, so the effect on the duration calculations is minor.

After the background, source, and model parameter selections have been made, every time bin in the selection is background subtracted, fitted using the model, and the resulting fitted spectrum is integrated over the appropriate energies to obtain a flux history. For comparison with the BATSE duration distribution, the energy range of 50–300 keV was chosen for this catalog. Errors for each integrated photon flux are derived using the covariance matrix for the fit, taking into account the uncertainties of each fitted model parameter. The resulting photon flux history (see, e.g., Figure 9) is summed over time, to produce a cumulative fluence plot, as seen in Figure 10. In this plot, the background-subtracted background intervals should, on the average, contribute zero to the total fluence, as seen in the left-hand portion of the figure, at times before the trigger time, and at the right-hand portion, well after the burst has concluded. In reality, depending on how well the seeded model fit parameters match with the fitted residuals, the flux histories in these two regions can exercise a random walk away from constant zero residual flux, as seen in Figure 11. Similar trends are present in the data in Figure 10 but they are small relative to the burst and hence not noticeable. In most cases, the random walk over the background accumulation does not exceed 1 photons cm\(^{-2}\), setting a practical limit of >2 photons cm\(^{-2}\) in total fluence for the duration analysis to be reliable. Treating this as a hard threshold would, however, bias the catalog against short bursts, which often have lower fluences, so results for weaker bursts are included herein but should be treated with caution.

To calculate the duration, the two “plateau” regions must be identified by hand (since every such plot is unique, this step cannot be automated), the average flux in each serving as the fiducial values against which the partial fluences at 5%, 25%,
and in turn, implies that the fluence levels at a given time, is also recorded for each of the time intervals that form the cumulative fluence. The variances of the two plateau regions can be used to derive several other important quantities. The variance measures our inability to precisely determine the zero and 100% levels. Once the variances in the flux histories are known, they are converted into uncertainties by first scaling by the desired flux level and taking the square root:

$$\text{error}_{\text{flux}} = \sqrt{(1-nn)^2 \text{Var}_0 + nn^2 \text{Var}_{100}}, \quad (A1)$$

where $nn \in [0.05, 0.25, 0.75, 0.95]$ indicates the various flux levels and $\text{Var}_0$ and $\text{Var}_{100}$ are the variances from the zero and 100% fluence level plateaus. For each of the four nominal fluence levels $nn$, $\text{error}_{\text{flux}}$ is added and subtracted, resulting in a projected uncertainty in time for each. The final uncertainty in $T_{90}$ is the root mean square of the corresponding uncertainties in $T_5$ and $T_95$ and similarly with the projected uncertainties in $T_{25}$ and $T_{75}$ for $T_{50}$. The start time, relative to the burst trigger time, is also recorded for each of the time intervals that form $T_{90}$ and $T_{50}$.

The flux history used in the calculation of the burst duration can be used to derive several other important quantities. The total fluence is calculated by differencing the zero and 100% integrated flux levels, as determined by the plateau selections. In each successive time bin the flux model from the fit is weighted by energy in erg, integrated over two energy bands, 50–300 keV and 10–1000 keV, and then added to the running total to produce the cumulative fluence. The variances of the two plateau regions are added together in quadrature to determine the uncertainty in the fluence. As with the $T_{90}$ calculation, it is the variance in the background regions (where the running sum should be zero) that determines the uncertainty of the zero level fluence (and similarly for the total fluence). Given the limited range of integration for the 50–300 keV fluences, the fact that we chose to use the COMP photon model makes little difference to the results; however, this may have a more significant effect on the 10–1000 keV fluences. The affected bursts are primarily the ones with significant emission in the BGOs, which are highlighted in Table 2. In any case, the fluences derived in this catalog are intended mostly as a ranking tool. Readers interested in more robust fluence estimates should consult the GBM spectroscopy catalog (Goldstein et al. 2012).

The deconvolved photon flux history is calculated by integrating the best-fit model for each time bin over the two energy bands described above. The peak flux is then the maximum value of the flux history between the lower and upper plateaus for the two energy bands, as well as for three different time intervals: 64, 256, and 1024 ms. As the native or default accumulation for the CTIME data post-trigger is 64 ms, there is only one possibility for binning, as long as care is taken to ensure that the peak flux interval occurred after $T = 0$, as is usually the case for long GRBs. For short GRBs, TTE is preferred, since much of the emission can occur pre-trigger, and TTE can be binned in 64 ms accumulations over its entirety. In order to calculate the 256 and 1024 ms peak fluxes, the available data are binned within a sliding window. CTIME pre-trigger accumulations are 256 ms by default, so only the post-trigger data need to be binned. In this case (CTIME), the peak flux is the maximum flux found either pre-trigger or in one of the sliding binning windows post-trigger. The time and value of the peak flux and its uncertainty are recorded, again calculated from the model fit, along with the uncertainties of the model parameters and the covariance matrix between the parameters, computed in the usual manner.

Note added in proof: A complete list of the IPN spacecraft that observed the bursts in Table 2 may be found at http://www.ssl.berkeley.edu/ipn3/masterl1.txt. An IPN supplement to the GBM catalog is in preparation (K. Hurley et al. 2012, in preparation).