Internal mass transport limitations inside Fischer–Tropsch catalysts due to the slow diffusion of reactants in the liquid-filled pores may significantly alter the selectivity and achievable productivity. In this work, diffusive restrictions for planar catalyst layers were investigated by mathematical modeling and simulation. A one-dimensional model utilizing empirical kinetics, incorporating transport pores as an additional pathway for mass transport and taking into account heat production, allows for calculation of catalyst efficiency and productivity towards C₅₊ products. As diffusional mass transport leads to strong concentration gradients that impair selectivity, an optimum layer thickness with maximum C₅₊ productivity can be found. Additional transport pores enhance the mass transport but reduce the amount of active phase, which requires a trade-off by optimizing the fraction of transport pores and layer thickness. For reference conditions, the catalyst layer with an ideal amount of transport pores and ideal thickness exhibits a productivity that is about 47% higher than that for the best layer without transport pores. This improvement requires transport pores with diameters not larger than about 60 μm. While the improvement potential significantly depends on the effective diffusivities, the effect of heat generation was found to be negligible.

### Introduction

Conversion of synthesis gas in low-temperature Fischer–Tropsch (FT) reaction yields a broad spectrum of hydrocarbons that can be used for various further applications such as the production of diesel fuels. An efficient process requires a high selectivity to long-chain hydrocarbons over the commonly used cobalt-based catalysts. Thus, products with high carbon numbers remain in liquid phase under the reaction conditions. Consequently, reactors for the low-temperature FT process must be suitable for the resulting three-phase system with fixed-bed reactors and slurry bubble columns being commercially established. Considering diffusion limitations, pressure drop and thermal behavior, the dimensions of catalyst pellets and tubes have to be optimized for conventional fixed-bed reactors. A further option is the use of micro-packed bed reactors, which are beneficial due to increased heat removal at moderate pressure drop. In reactors of this type, highly active catalysts exhibit excellent performance, but depending on the catalyst particle size, one can either improve the pressure drop or the mass transport. Microstructured reactors with a wall coating of a catalyst offer an extraordinarily low pressure drop and high catalyst efficiency but suffer from low catalyst inventory and therefore low reactor productivity. The low catalyst inventory results from diffusion limitations that restrict the thickness of the catalyst coating. Diffusion inside the FT catalyst is slow due to the liquid products that remain inside the pores of the catalyst during the reaction. Moreover, these severe internal mass transport effects do not only retard the reaction but also hamper selectivity.

Different approaches have been made to deal with the restrictions in catalyst size caused by limited diffusion length. To achieve a low diffusion length, the use of eggshell catalysts or structured packings and open foams as catalyst supports was proposed. This reduces the negative impact of internal diffusion limitations, results in a low pressure drop and, especially for foams and packings, also good heat removal. Nonetheless, these approaches suffer from low catalyst inventory and do not directly enhance internal mass transport. For improvement of the diffusive transport inside the catalyst, a bimodal pore structure can be used. Compared experimental results from FT synthesis with a mathematical model, indicating that a bimodal catalyst gives rise to high activity at low diffusion resistance.
pores confine the cobalt species to nanometer-sized particles with high specific surface areas. On the other hand, diffusion becomes very slow in liquid-filled catalyst support pores of small size, as shown by Preising and Enke for porous glasses. This can be attributed to restricted diffusion as a result of a solute with a critical molecular diameter in the range of the pore diameter. This pore size effect is similar to the well-known Knudsen diffusion regime during gas-phase mass transport. Although a comprehensive, exact matching correlation for the description of these effects could not yet be developed, the work of the abovementioned authors emphasizes that nanometer-sized pores additionally hinder diffusion in the liquid phase. Therefore, optimized catalyst designs require larger pores to improve the accessibility of the catalyst for the reactants in addition to smaller pores providing high surface area and activity.

Different approaches to describe the effect of pore structures on effective diffusion are reported in the literature. From general analytical solutions of simplified models for bidisperse systems or fractal pore models to numerical approaches regarding three-dimensional pore networks, catalyst efficiency depends on effective diffusivity and thus on structural parameters like tortuosity or pore connectivity. Although these parameters can be measured, the validity, especially for tortuosity, under the reaction conditions is not necessarily given. A three-dimensional multiscale simulation approach aims to avoid these uncertainties by starting with a nanometer scale model to evaluate the effective diffusivities for a mesoporous structure, which can then be used for further simulation of effective diffusivities and reaction rates on the micrometer scale. In the last step, these results are then applied to evaluate a reactor model. For CO oxidation, this approach by Pereira et al. leads to a good agreement between the experimental results and the simulation. But even here, the validity of effective diffusivities depends on the reliability of the nanoscopic structure. An optimization by Gheorghiu and Coppens of two-dimensional bimodal pore networks for a simple first-order gas-phase reaction illustrates the trade-off between ease of access and productivity, which depends on the pore fraction and pore size. It is shown that transport pores are especially interesting if diffusion in micro- and mesopores is slow due to Knudsen diffusion. Further systematic evaluation by Wang et al. and Johannessen et al. has revealed that uniform transport pore size and fraction can yield almost the same improvement as an optimized distribution of transport pores. These results obtained for simple first-order kinetics were applied to calculate the optimal porous structure of a catalyst layer for methane reforming, indicating a considerable increase in activity for a commercial catalyst.

In the present work, a simplified mathematical model is used to describe the effects of reactant diffusion in catalyst layers with small and large pores on the activity and selectivity of the FT reaction. Using the approach of Vervloet et al. for description of the chain growth probability as a function of reactant concentration and temperature, we have already shown that it is possible to find an ideal layer thickness, where the yield with respect to the desired products reaches a maximum. In the present paper, the optimization potential is analyzed in detail with respect to the effective diffusivities in the catalyst and the transport pore phase by adjusting the transport pore fraction. Furthermore, the relevance of temperature gradients and the maximum allowable transport pore diameter are evaluated. In other words, by applying realistic, experimentally based kinetics and transport parameters to optimization of the pore structure, we demonstrate how and up to what extent transport pores can enhance the productivity in Fischer–Tropsch synthesis.

**Model description**

The general idea is that transport pores inside a catalyst provide a bypass for diffusive transport of reactants inside the catalyst and thus improve the effective mass transport. However, since the volume fraction of transport pores does not contribute to the reaction, an increased transport pore fraction improves mass transport at the expense of a reduced amount of the catalytically active phase. This emphasizes the necessity of a compromise that has to take into account the reduced volume fraction of the active phase and the benefit of an improved mass transport.

The developed model is based on a planar geometry of a catalyst coated onto the wall of a microchannel reactor, but can also be applied to wascoated monoliths and eggshell catalysts, as long as the active catalyst layer remains thin. The catalyst layer is assumed to consist of a mesoporous active catalyst and cylindrical macropores denoted as transport pores.

Above the catalyst layer lies the free cross-section area of the channel. In order to investigate the effects of the diffusion in the catalyst and the transport pores only, axial convection and external mass transport are neglected. This leads to the assumption of constant reactant concentrations at the gas–solid boundary and simplifies the model to only one dimension (Fig. 1). Although this reduction to a differential reactor does not lead to results which could directly be compared to experimental data, it allows results with general significance to be derived, which are also applicable to different reactor concepts. For the solution of the model equations, gPROMS® ModelBuilder 4.0.0 of Process Systems Enterprise Ltd. was used. The transport equations were distributed over 50 elements using a second-order “orthogonal collocation on finite elements method” (OCFEM) for discretisation.

**Transport equations**

The reaction takes place on the active surface inside the catalyst. Therefore, the reactants, carbon monoxide and hydrogen, have to be transported into the catalyst layer. This diffusive transport occurs as a result of a concentration gradient between the gas–solid boundary and the depth of the catalyst. Since two types of pores concurrently contribute to diffusive mass transport but the reaction occurs on the active
surface only, the catalyst phase and the transport pore phase must be distinguished within the catalyst layer. The volume fraction of the transport pores is defined by $\varepsilon_{TP}$. This virtual separation of the transport phase and the active phase is a main idea of the simulation approach (Fig. 1(a)).

The representative volume element for the mass balance is the whole catalyst layer (Fig. 1(a)) comprising the active and transport phases. Mass transport is described according to eqn (1) with a term considering the effective diffusion inside the catalyst and the transport pore phase, and a reaction term. As the transport pores are inert, the active catalyst fraction $(1 - \varepsilon_{TP})$ has to be considered for the reaction term. The diffusion term simplifies the original two-dimensional problem to a one-dimensional equation, accounting for the different effective diffusivities inside the catalyst and the transport pores. This effective one-dimensional model is valid as long as the transport pores are sufficiently small, thereby leading to the same concentration profiles in the catalyst and transport pores.47

$$
\left(1 - \varepsilon_{TP}\right)\varepsilon_{Cat} \frac{D \varepsilon_{Cat}}{\varepsilon_{TP}} \frac{d^2}{dx^2} \left(c(x)\right) + (1 - \varepsilon_{TP}) \nu (x) r_{CO}(x) = 0 \quad (1)
$$

The effective diffusion is assumed to obey Fick’s law which is modified with porosity and tortuosity. It has to be mentioned that the reference volume for the catalyst porosity is the catalyst phase, while that for the transport pore fraction is the whole catalyst layer. Thus, the sum of the transport pore fraction and catalyst porosity is not unity. The assumed reference values for porosity and the tortuosities are listed in Table 1. As the transport pores are cylindrical and parallel to the direction of the diffusion, the tortuosity of the transport pores equals one. This is an ideal assumption, since the transport pores of real catalyst layers will not easily exhibit such a simplified best case geometry. The reference tortuosity of the catalyst is assigned a value of 3, which is in the typical range.29,50 The supposed reaction kinetics in eqn (1) are dependent on the concentration of carbon monoxide and hydrogen; consequently, only these species are relevant for mass transport and needed to be considered for the boundary conditions (eqn (2) and (3)).

$$
c|_{x=0} = \frac{p}{f v} \quad (2)
$$

$$
\frac{\partial c}{\partial x} |_{x=d} = 0 \quad (3)
$$

Since the FT reaction is quite exothermic, the removal of heat may have a significant impact. Therefore, the heat balance was taken into account. The corresponding differential eqn (4) neglects the influence of a varying transport pore fraction on the effective thermal conductivity and incorporates the transport pore fraction only via the residual amount of catalyst $(1 - \varepsilon_{TP})$, which scales the heat source term.

$$
\lambda_{eff} \frac{d^2}{dx^2} \left(T(x)\right) + (1 - \varepsilon_{TP}) r_{CO}(x) \Delta_e H = 0 \quad (4)
$$

Similar to mass transport, the boundary conditions assume that all heat is removed through an isothermal wall and the heat flux to the gas boundary is zero (eqn (5) and (6)).

$$
T|_{x=t} = T_0 \quad (5)
$$

$$
\frac{\partial T}{\partial x} |_{x=0} = 0 \quad (6)
$$

| Parameter | Value | Unit |
|-----------|-------|------|
| $\varepsilon_{Cat}$ | 0.4 | $m_{pore}^3 m_{Cat}^{-3}$ |
| $\tau_{Cat}$ | 3 | — |
| $\tau_{TP}$ | 1 | — |
| $\lambda_{eff}$ | 0.1 | W m$^{-3}$ K$^{-1}$ |
For the reaction enthalpy, a value of 170 kJ mol\(^{-1}\) was chosen, which is a quite high estimate compared to the literature.\(^{5,6,11,51}\) For the effective thermal conductivity, a value was chosen, which is significantly lower than the conductivities of the solid catalyst phase,\(^{48,52}\) the liquid phase,\(^{53}\) the gas phase in the Fischer–Tropsch reaction\(^{54}\) or the estimated effective bed conductivity.\(^{55}\) Hence, the chosen approach should give rise to a worst case assessment of possible heat transport effects.

### Kinetics and selectivity

For the description of the CO consumption rate, an approach by Yates and Satterfield\(^{56}\) is used (eqn (7)), where a temperature dependence analogous to the Arrhenius equation as proposed by Maretto and Krishna\(^{57}\) is utilized (eqn (8) and (9)).

\[
\frac{\nu_{\text{CO}}}{P_{\text{CO}}} = \frac{a}{1+b} \left( P_{\text{H}_2} P_{\text{CO}} \right)^{\frac{a}{b}}
\]

\[
a = a_0 \exp \left( \frac{E_{\text{AA}}}{R} \left( \frac{1}{493.15 K} - \frac{1}{T} \right) \right)
\]

\[
b = b_0 \exp \left( \frac{E_{\text{AB}}}{R} \left( \frac{1}{493.15 K} - \frac{1}{T} \right) \right)
\]

Assuming paraffins to be the only hydrocarbon products (eqn (10)), the chain growth probability allows the calculation of the stoichiometric coefficients for all species participating in the reaction.

\[
\nu_{\text{CO}} \text{CO} + \nu_{\text{H}_2} \text{H}_2 \rightarrow \nu_{\text{H}_2O} \text{H}_2 \text{O} + \nu_{\text{C}_n\text{H}_{2n+2}} \text{C}_n\text{H}_{2n+2}
\]

The relationship developed by Vervloet \textit{et al.}\(^{18}\) describes the chain growth probability as a function of reactant concentration and temperature (eqn (11)). Constants for the reaction rate and chain growth probability are specified in Table 2. The stoichiometric coefficients of carbon monoxide and water are independent of the hydrocarbon chain length (eqn (12) and (13)). However, this is not the case for the consumption of hydrogen, for which the stoichiometric coefficient is accessible via eqn (14).\(^{18}\) For a chain growth probability of zero, all carbon monoxide is converted to methane, thus leading to a stoichiometric coefficient of three for hydrogen, whereas for a chain growth probability close to unity, long-chain hydrocarbons are formed almost exclusively, hence resulting in a hydrogen coefficient of virtually two. The stoichiometric coefficients of the hydrocarbons follow the ASF distribution according to eqn (15).

\[
\alpha = \frac{1}{1 + k_\alpha \left( \frac{c_{\text{H}_2}}{c_{\text{CO}}} \right)^{\beta} \exp \left( \frac{\Delta F_\alpha}{R} \left( \frac{1}{T_{\text{ref}}} - \frac{1}{T} \right) \right)}
\]

\[\nu_{\text{CO}} = -1\]

\[\nu_{\text{H}_2O} = 1\]

\[\nu_{\text{H}_2} = -(3 - \alpha)\]

\[\nu_{\text{C}_n\text{H}_{2n+2}} = \alpha^{n-1}(1 - \alpha)^2\]

The expressions for the chain growth probability and the overall reaction rate are functions of temperature and the concentrations of hydrogen and carbon monoxide only. As the stoichiometric coefficients solely depend on the chain growth probability, it is sufficient to solve the mass transport eqn (1) and (2) exclusively for hydrogen and carbon monoxide.

### Objective of the simulation

Simulation of the reaction–diffusion equations computes coupled variables as spatial distribution over the catalyst layer thickness. These variables are the hydrogen and carbon monoxide concentrations, reaction rate, temperature, chain growth probability and the resulting stoichiometry. Computation utilizes transport equations, kinetic equations, material properties and parameters to estimate the values for all distributed variables. For evaluation and comparison of the overall reactivity of different catalyst layers, the catalyst efficiency is a measure, which compares the apparent reaction rate with the surface reaction rate (eqn (16)). Since catalyst layers with a certain amount of transport pores exhibit a lowered amount of the catalytically active phase, the efficiency has to be corrected by the residual fraction of the catalytically active phase (eqn (17)) leading to the layer efficiency. For this new quantity, the volume of the catalyst layer serves as the reference volume, which facilitates a reliable comparison of different layer productivities.

\[
\eta_{\text{C}_n\text{H}_{2n+2}} = \frac{1}{\eta_{\text{C}_n\text{H}_{2n+2}}} \int_0^1 \frac{r_{\text{CO}}(x)}{r_{\text{CO}}(x=0)} \, dx
\]

\[\eta_{\text{C}_n\text{H}_{2n+2}} = \frac{1}{\eta_{\text{C}_n\text{H}_{2n+2}}} \int_0^1 \frac{r_{\text{CO}}(x)}{r_{\text{CO}}(x=0)} \, dx\]
\[ \eta_{co} = (1 - \varepsilon_{TP}) \int_{0}^{r_{CO}(x)} f \, r_{CO}(x) \, dx \] (17)

To evaluate the fraction of the desired products, which are assumed to consist of hydrocarbons with carbon numbers of 5 or higher, the C5+ selectivity on a carbon mass basis is calculated as a function of spatial stoichiometric coefficients, which directly depend on the chain growth probability (eqn (18) and (19)). For further classification of the values obtained with this equation, the methane selectivity is calculated according to eqn (20). From these equations, the effective selectivities are accessible via integration over the whole layer thickness and averaging with the reaction rates (eqn (21) and (22)).

\[ S_{CO}(x) = n \alpha(x)^{n-1} (1 - \alpha(x))^2 \] (18)

\[ S_{C5+}(x) = 1 - \sum_{n=5}^{10} S_{CO}(x) = (5 - 4 \alpha(x)) \alpha(x)^4 \] (19)

\[ S_{CL}(x) = (1 - \alpha(x))^2 \] (20)

\[ S_{C5+eff} = \frac{\int_{0}^{r_{CO}(x)} S_{C5+}(x) \, dx}{\int_{0}^{r_{CO}(x)} \, dx} \] (21)

\[ S_{CLeff} = \frac{\int_{0}^{r_{CO}(x)} S_{CL}(x) \, dx}{\int_{0}^{r_{CO}(x)} \, dx} \] (22)

Typically, the space time yield, STY, is used to compare the productivity of different reaction systems. In this particular case—due to the flat geometry of the catalyst layers and fixed gas composition—gradients in concentration will only appear in the direction of the layer thickness. This means that the absolute overall productivity scales nonlinearly with the thickness but is proportional to the geometric surface area of the catalyst layer. Thus, an area-specific quantity denoted by areal time yield, ATY, is a better reference and therefore used as an objective function for evaluation of the transport pore effects on mass transport. ATY can be calculated from the distributed variables of the reaction rate and C5+ selectivity (eqn (23)) and is also equal to the molar flux of C5+, carbon species leaving the layer.

\[ ATY = (1 - \varepsilon_{TP}) \int_{0}^{r_{CO}(x)} S_{C5+}(x) \, dx \] (23)

The focus of this investigation is to find the maximum productivity of an optimally designed catalyst layer. For wall-coated micro-reactors as well as eggshell catalysts and washcoated monoliths, catalyst inventory rises with increasing layer thickness. This implies the demand to fill as much catalyst as possible into the reactor by increasing the layer thickness, but simultaneously avoiding diffusion limitations. Hence, one objective is to find the layer thickness that yields the highest productivity. The other objective is to evaluate the influence of the transport pore fraction on the achievable productivity. Therefore, two degrees of freedom are available, namely the catalyst layer thickness and the transport pore fraction. The ATY is advantageous for the desired evaluation since it allows the optimization towards both parameters to be performed, while evaluating the catalyst efficiency or STY only is insufficient for this task (see detailed discussion in the ESI†). However, for comparison of layers with the same thickness, only the efficiency and selectivity have to be considered as the thickness and surface reaction rate remain constant.

**Physical properties**

Computation of the mass transport equation requires certain properties of hydrogen and carbon monoxide. The gas phase is supposed to be ideal and contain only hydrogen and carbon monoxide in a fixed H2/CO ratio of 2 at a temperature of 493.15 K and a pressure of 21 bar. Thus, the concentration and diffusivity of the reactants in the liquid phase only depend on the liquid solvent. Although the reaction produces a broad distribution of hydrocarbons, the liquid phase is assumed to consist of only one species for the estimation of liquid solubility and diffusivity. For a paraffin with a carbon number of 28, the liquid molar volume and the Henry constants for hydrogen and carbon monoxide are calculated as reported by Marano and Holder.58 Diffusivities of hydrogen and carbon monoxide are calculated in accordance with an approach by Erkey et al.59 The resulting properties for the reaction conditions are listed in Table 3.

| Parameter | Value | Unit |
|-----------|-------|------|
| \( T_0 \) | 493.15 | K |
| \( p \) | 21 | bar |
| \( H_2/CO \) | 2.0 | mol mol\(^{-1}\) |
| \( c_{H_2} \) | 52.46 | mol m\(^{-3}\) |
| \( c_{CO} \) | 33.07 | mol m\(^{-3}\) |
| \( H_{H_2} \) | 458.6 | bar |
| \( H_{CO} \) | 363.8 | bar |
| \( v_L \) | 5.818 \times 10^{-3} | m\(^3\) mol\(^{-1}\) |
| \( D_{H_2} \) | 14.30 \times 10^{-9} | m\(^2\) s\(^{-1}\) |
| \( D_{CO} \) | 36.05 \times 10^{-9} | m\(^2\) s\(^{-1}\) |
Simulation results

Intrinsic reaction rates

Fig. 2 displays the consumption rate of CO, the C₅⁺ selectivity and the resulting formation rate of C₅⁺ hydrocarbons as functions of reactant gas composition.

Since no mass transport effects are considered and only the kinetic expressions are relevant, the rates are denoted as intrinsic. At small H₂/CO ratios, where the concentration of carbon monoxide is high, the C₅⁺ selectivity reaches maximum values. However, the inhibiting effect of carbon monoxide in the kinetic expression (eqn (7)) limits the achievable reaction rate. On the contrary, a low carbon monoxide concentration promotes the overall reaction rate but confines the reaction rate. These opposite trends lead to a maximum of the formation rate of the desired products at a distinct value for the H₂/CO ratio of slightly above 3. Although diffusional effects will lower the concentration of both reactants in the catalyst layer, a shift in the reactant ratio will also appear, especially as CO and H₂ exhibit different diffusivities.

Reaction–diffusion effects – limitation of layer thickness

Fig. 3 illustrates the effect of coupled reaction and diffusion on the spatial concentration profiles inside catalyst layers of different thicknesses without any transport pores. The concentration is normalized to the concentration of hydrogen at the gas–solid boundary, which is given in Table 3. As this is the highest observed concentration, the normalized concentration profiles in Fig. 3 range between the values of zero and unity. Additionally, the resulting chain growth probability is plotted in the same figure. For the smallest thickness of 10 μm, gradients in concentration are almost absent and the normalized concentrations remain constant at a value of unity for hydrogen and about 0.63 for carbon monoxide (Fig. 3, top left). This increased carbon monoxide concentration results from the lower Henry coefficient of CO compared to H₂ (Table 3), which leads to a better solubility of CO (eqn (2)). Thus, the H₂/CO ratio inside the liquid is 1.59 for a value of 2 in the gas phase. The chain growth probability also stays at a level of about 0.89 due to the absence of gradients.

With increasing diffusion length, concentration gradients start to form as can be seen for a layer thickness of 100 μm (Fig. 3, top right). Although the relative concentrations of hydrogen and carbon monoxide drop to 0.85 and 0.45, respectively, the influence on the chain growth probability is almost negligible. However, after a further increase of the layer thickness to 150 μm (Fig. 3, bottom left), a severe drop of reactant concentrations occurs. At the wall boundary, the normalized concentration of carbon monoxide is decreased to 0.05, whereas the concentration of hydrogen still remains above 40% of the original value. As a result, the H₂/CO ratio reaches a value of about 8 at this position, thereby diminishing the chain growth probability to 0.30. For an even larger layer thickness of 300 μm (Fig. 3, bottom right), the concentration of carbon monoxide reaches a value of virtually zero directly behind the middle of the layer thickness and from this point on, the concentration of hydrogen persists at about one third of its initial value. This almost total absence of CO causes a drastic increase of the H₂/CO ratio which leads to a chain growth probability of effectively zero. In particular, the last two cases illustrate that the better diffusivity of hydrogen compared to carbon monoxide leads to an increased H₂/CO ratio, which negatively affects the chain growth probability. Neither the better solubility of carbon monoxide nor the more than twice as high consumption rate of hydrogen compensates this effect.

In accordance with Fig. 3, the CO consumption rates as well as the methane and C₅⁺ selectivity profiles are depicted in Fig. 4 for four catalyst layers of different thicknesses. The reaction rate is normalized to the intrinsic value at the surface boundary. Since no substantial concentration gradients occur at a layer thickness of 10 μm, the reaction rates and selectivities exhibit a constant value (Fig. 4, top left). For the 100 μm layer, the minor concentration gradients with a slight increase in the local H₂/CO ratio result in an only minor change of methane selectivity from 0.01 at the surface to a value of 0.02 at the wall boundary. Nonetheless, the C₅⁺ selectivity is already quite clearly influenced and decreases from 0.90 to 0.84. However, due to the increased H₂/CO ratio, the inhibiting effect of CO on the reaction rate is reduced and the rate rises to 12% above the intrinsic value (Fig. 4, top right). These contrary trends for the C₅⁺ selectivity and reaction rate become more relevant at a layer thickness of 150 μm. Although the concentration of carbon monoxide is only marginally above zero in proximity to the wall, a normalized rate of 1.87 is calculated at this position (Fig. 4, bottom left). Because of the increased local H₂/CO ratio, the C₅⁺ selectivity is now severely diminished and at the wall, only 3% of the converted carbon monoxide participates in the production of the desired C₅⁺ products, whereas almost 50% is converted to methane. At a layer thickness of 300 μm, the concentration profile of carbon monoxide...
reaches zero at approximately half of the thickness and consequently the reaction is stopped. However, the profile of the reaction rate exhibits a strong increase up to this point which drastically drops when almost all carbon monoxide is...
consumed. This leads to a distinctive maximum in the reaction rate profile for catalyst layers with strong mass transfer limitations. At the point of the maximum reaction rate, the C_{5+} selectivity reaches zero and the methane selectivity considerably increases towards unity. This means that an effective contribution to the production of C_{5+} hydrocarbons is limited to the first third of the catalyst layer. Any additional catalyst thickness can be denoted as dead volume since the formation rate of C_{5+} is practically zero. The profiles of the selectivity and reaction rate reveal that increasing the layer thickness has a negative effect on the selectivity to the desired products and a more complex effect on the efficiency as the reaction rate displays a maximum. Especially, the last case with 300 \mu m thickness is in good agreement with the previously reported results by Vervloet et al.\cite{18} The concentration profiles are quite similar, but since the reaction conditions differ slightly and a spherical model instead of a planar geometry was used, the dimensionless rate profile exhibits a higher maximum with a normalized rate above 2.5 in the work of Vervloet et al.

For comparison of the productivity of catalyst layers with different thicknesses, the effective selectivities, efficiencies and ATY are shown in Fig. 5. On the basis of the ATY, the productivity can be divided into three different regions. In the first one ranging up to a layer thickness of approximately 100 \mu m, concentration gradients are not relevant. This means that the ATY is directly proportional to the layer thickness as the C_{5+} selectivity and reaction rate remain close to the intrinsic values, which transforms the integral in eqn (23) into the expression ATY = (1 - \epsilon_{TP}) r_{CO}S_{C_{5}}, L. In the next region for thicknesses ranging from ca. 100 \mu m to 160 \mu m, concentration gradients develop leading to partially opposing tendencies for efficiency and selectivity. As the increase in layer efficiency compensates for the decreased selectivity only to some extent, a maximum for the ATY at a layer thickness of 139 \mu m occurs. Although from this point on the layer efficiency still increases to its maximum of 1.45 at 151 \mu m, a higher efficiency does not lead to a higher ATY as the drop in C_{5+} selectivity is too severe. Above a layer thickness of about 160 \mu m, the C_{5+} selectivity and ATY remain constant. This marks the third and last region from which any further increase in layer thickness would only result in an additional amount of the active phase that does not contribute to the formation of the desired products. Fig. 5 displays a characteristic profile of the ATY that allows a certain layer thickness which produces the highest achievable amount of desired products for catalyst layers without transport pores to be found. Hence, it provides a reference for evaluating the effect of transport pores on mass transport enhancement and improving the catalyst layer productivity. This reference is denoted as optimum for “dense” layers.

**Mass transfer improvement by transport porosity optimization**

The results of a systematic variation of transport porosity for the case with a layer thickness of 300 \mu m and a transport pore tortuosity of one are shown in Fig. 6. Transport porosity is varied from the lower theoretical boundary of zero to a maximum value of 0.99. The effective values for layer efficiency, selectivities and ATY are plotted as functions of transport porosity. A catalyst layer of 300 \mu m exhibits a strong influence of concentration gradients if no transport pores are present. As a result, the efficiency is limited to about 0.74, the C_{5+} selectivity is fairly low with a value of about 0.40 and the methane selectivity of 0.32 is quite high. With increasing transport porosity, the selectivities to C_{5+} hydrocarbons and methane remain unaffected up to a porosity of 25%. Only the layer efficiency rises, which can be attributed to an improved accessibility of deeper catalyst layer fractions.

As long as an enhanced access of reactants, especially carbon monoxide, does not prevent a concentration drop towards zero, a high local H_{2}/CO ratio still prevails and results in constant selectivities. With further increasing

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**Fig. 5** Integral values for catalyst efficiency, C_{5+} selectivity and areal time yield of catalyst layers without transport pores as a function of thickness.

**Fig. 6** Layer efficiency, areal time yield, C_{5+}, and methane selectivity of a 300 \mu m thick layer as functions of transport pore fraction.
porosity, the negative effect of high local H₂/CO ratios disappears and selectivities rapidly approach their intrinsic values. On the other hand, the positive effect of high H₂/CO ratios on the reaction rate vanishes as the diffusive transport is enhanced. Together with a reduced amount of the active phase, this leads to decreasing layer efficiency. At high fractions of the transport pores, diffusive transport is fast enough for an almost intrinsic behavior throughout the catalyst layer. Thus, the efficiency drops to a value proportional to the residual amount of the catalyst phase. As a consequence, the ATY is only slightly enhanced by increased efficiency up to a porosity of 25% but fairly benefits from an improved selectivity passing a maximum at a porosity of 33% before it declines in accordance with the layer efficiency at higher porosities. At maximum, the achievable productivity is more than twice as high as for the case without transport pores. Fig. 6 nicely demonstrates the positive effect of the optimal fraction of transport pores. However, one should keep in mind that not only the transport pore fraction but also the layer thickness is required to be optimized in order to compare only the best points to determine a reliable improvement. For this purpose, the results of a single layer optimization of transport porosity, as previously discussed, are condensed in Fig. 7 for catalyst layers ranging from 0 to 500 μm.

For each layer thickness, the maximum productivity represented by ATY, the corresponding layer efficiency, the C₅⁺ selectivity and the required transport porosity are collected. Thus, Fig. 7 represents the optimization results for each single layer thickness. For thin layers, diffusion limitations are absent and additional transport pores only have a negative effect, as they reduce the fraction of the active phase. Hence, up to a layer thickness of 135 μm, catalytic layers without transport pores yield the highest ATY and the plots are equivalent to those in Fig. 5. However, from this thickness and higher, the introduction of transport pores enhances the achievable productivity. The C₅⁺ selectivity remains almost constant at a value between 0.9 and 0.7 and a distinct maximum of layer efficiency is extenuated leading to a typical hyperbolic decrease with increasing layer thickness. The necessary fraction of transport pores rises non-linearly from zero to above 60% at a layer thickness of 500 μm. The increasing transport porosity enables the values of ATY to increase for high thicknesses significantly above the level of non-optimized layers. Nonetheless, it is not possible to continue the linear increase of ATY by optimization of transport porosity and even for this ideal case, a maximum of ATY over the layer thickness appears. The maximum achievable ATY at the ideal transport pore fraction of 0.43 and at an ideal thickness of 356 μm is about 47% higher than the maximum for “dense” layers without transport pores. The thickness where the maximum occurs is significantly shifted to higher values, but since transport pores are required to obtain these higher thicknesses, the amount of the catalyst is not increased proportionally. For the optimum case, the layer thickness at maximum ATY corrected with the residual fraction of the catalyst phase of 0.57 exhibits a value of 204 μm. This is 47% higher than that for the dense case and underlines that the enhanced ATY is directly related to an increased amount of the catalyst. Moreover, transport pores are necessary at this elevated diffusion length to preserve a concentration profile with minor gradients. These minor gradients allow a relatively high C₅⁺ selectivity of 0.73 to be maintained, which is almost identical to the C₅⁺ selectivity for the optimum of the dense case (Fig. 5).

At first sight, an enhancement of only 47% appears to be relatively low compared to improvement through addition of transport pores reported in the literature. However, it has to be taken into account that only the maxima of ATY for the dense and ideal cases are compared. This comes along with different values for transport porosity and layer thickness. If one compares the productivity at a certain layer thickness only, significantly higher values will be obtainable, e.g. for a layer thickness of 500 μm, increases up to 97% are possible. This emphasizes the necessity to compare only the maxima for different cases and to include not only the transport pore fraction but also the catalyst layer thickness during optimization.

### Effect of tortuosity

Since the model assumes straight, cylindrical pores perpendicular to the layer surface, the tortuosity of transport pores equals unity. For real catalyst structures, this ideal value is not necessarily reached and diffusivity inside transport pores might be slower due to higher values of transport pore tortuosity. Furthermore, the tortuosity of the mesoporous catalyst can be lower or higher than the assumed value for the reference case (τₚ = 3). For the dense and the optimized ideal layers, the achievable maxima of ATY as a function of catalyst and transport pore tortuosity are shown in Fig. 8. Without transport pores, the ATY considerably decreases with increasing catalyst tortuosity. For the best points with ideal transport pores (τₚ = 1), the overall highest productivity can be obtained at a catalyst tortuosity of one. The decline with
increasing catalyst tortuosity is considerably attenuated almost leading to a constant productivity at high values of $\tau_{\text{Cat}}$. With increasing transport pore tortuosity, the diffusion in the additional pores slows down and thus the achievable ATY drops. For the points where the catalyst tortuosity is equal to or even lower than the transport pore tortuosity, no or only a negligible improvement can be observed. This clearly emphasizes the necessity of transport pores with high effective diffusivities in order to facilitate mass transport and obtain a benefit. In Fig. 9, the required transport pore fraction as a function of catalyst and transport pore tortuosity is plotted. With increasing catalyst tortuosity, the required porosity rises, indicating that especially at slow diffusion inside the catalyst, introduction of an inert transport pore volume gives rise to an improvement. The similar shape of the profiles in Fig. 8 and 9 indicates that only the difference of the diffusivity inside the catalyst and the transport pores is relevant for the improvement and the required transport pore fraction.

Thus, the achievable enhancement by insertion of transport pores is directly visible as a function of the ratio of catalyst to transport pore tortuosity (Fig. 10, calculated with the data from Fig. 8 and 9). It is evident that the diffusion in the transport pores has to be faster than the diffusion inside the catalyst to yield an effect for ATY. Furthermore, the improvement potential increases with rising tortuosity ratio. Interestingly, the required transport porosity rapidly approaches a value close to 0.5 with increasing tortuosity ratio. This is a very similar result to that reported by Coppens and coworkers, who found a maximum ideal porosity of macro pores of 0.5 for layers with the highest investigated Thiele moduli.

Maximum allowable transport pore diameter

To justify the reliability of an effective 1D model, transport pores have to be sufficiently small in order to avoid any additional transport limitations orthogonal to the main diffusion direction. Johannessen et al. have derived a criterion for the necessary size of transport channels from 2D simulations for first-order kinetics and applied this to cylindrical transport pores in a hexagonal packing. For the present work, the criterion can be written as eqn (24), assuming CO as the limiting species due to its lower diffusivity and the approximation of the first-order kinetics using the reaction rate and concentration at the gas–solid boundary.

$$d_{\text{TP}} \leq 0.2 \sqrt{\frac{D_{\text{Cat}} c_{\text{CO}}(x = 0)}{r(x = 0)}} \sqrt{\frac{2\sqrt{3}}{\pi} e_{\text{TP}}} \left(1 - \frac{2\sqrt{3}}{\pi} e_{\text{TP}}\right)^{-1}$$

This criterion basically determines the maximal remaining wall thickness between the transport pores, in order to avoid
concentration gradients inside the walls. Therefore, the pore size is independent of the layer thickness and only depends on the effective diffusion and the reaction inside the catalyst phase as well as the fraction of transport pores.

For catalyst tortuosities of 1, 3 and 10, the resulting wall thicknesses are about 49 μm, 28 μm and 15 μm, which are considerably smaller than the layer thickness at which concentration gradients appear (Fig. 5). The transport pore diameters calculated by this criterion for the three different catalyst tortuosities are shown in Fig. 11. Even for a very low transport pore fraction of 0.01 and the lowest considered catalyst tortuosity of τ_{Cat} = 10, the maximum allowable transport pore diameter is still about 1.8 μm. Thus, transport pores with diameters of 1–2 μm can be used over the entire range of layer thicknesses and transport pore fractions without causing additional mass transfer resistances. However, the required transport pore fraction for the optimal point is, in most cases, in the range of 0.4 to 0.5 (Fig. 10). Here, the necessary transport pore diameter can be between 30 and 140 μm. This shows that remarkably large pores can be utilized to achieve an overall improvement of FTS.

**Temperature gradients**

Although the main focus of this work is the enhancement of mass transport effects, we also investigated the possible appearance of temperature gradients in the catalyst layers. The maximum temperature deviations from the wall temperature occur at the gas–solid boundary as no heat flux to the gas phase was assumed. Fig. 12 displays the temperature differences at this point as a function of layer thickness. Obviously, the temperature difference rises with increasing layer thickness. However, even for a layer of 500 μm, the calculated temperature rise is less than about 0.5 K. Consequently, any influence of the temperature distribution in the catalyst layer can be neglected.

**Conclusions**

Diffusion phenomena inside Fischer–Tropsch catalysts are likely to play a vital role in the apparent reaction rate and selectivity to long-chain hydrocarbons and may limit the achievable productivity. The insertion of transport pores is an interesting option to enhance catalyst productivity. For the investigation of this concept, an effective one-dimensional reaction–diffusion model using empirical kinetic equations was adapted to include transport pores as the inert phase. For a planar geometry of the catalyst, the quantity ATY was defined and used as the objective function for productivity optimization. During optimization, it became evident that the catalyst thickness and transport pore fraction have to be adjusted simultaneously in order to obtain the maximum benefit in productivity. For the investigated parameter range, an improvement of up to 150% was found, strongly depending on the effective diffusivities inside the transport pores and the catalyst phase. The best results are achieved if the diffusion in the transport pores is significantly faster than that inside the catalyst. In this range, a relatively constant transport pore fraction of 0.4 to 0.5 was found to be ideal. For the reference case, the productivity could be improved by about 47% with a transport pore fraction of 43%. To achieve this enhancement, the layer thickness has to be increased from 134 μm to 346 μm. The maximum allowable transport pore diameter at this point is 62 μm. As for varying tortuosities, different transport pore fractions were found to be ideal, and the limit for the transport pore diameter also changed. In most cases, transport pore diameters between 30 and 140 μm can be employed without any additional mass transfer resistance between the transport pore and the catalyst. The calculated temperature gradients in the catalyst layer are fairly low. Thus, mass transport is the dominant phenomenon determining the rate and selectivity of FTS catalysts.

The simple model described in the present contribution illustrates that variation of structural parameters offers...
considerable optimization potential for FTS catalysts. These parameters can be easily altered with common techniques of catalyst preparation such as addition of pore-forming agents or templates. These investigations are currently carried out by our group and will be presented in a forthcoming publication.

**Notation**

**Latin**

- $a$: Reaction rate constant, mol s$^{-1}$ bar$^{-2}$
- $a_0$: Frequency factor for reaction rate constant, mol s$^{-1}$ bar$^{-2}$
- ATY: Areal time yield of desired products, mol s$^{-1}$ m$^{-2}$
- $b$: Absorption constant, bar$^{-1}$
- $b_0$: Frequency factor for absorption constant, bar$^{-1}$
- $c_{i,\text{Cat}}$: Concentration in liquid-filled pores of the catalyst phase of species $i$ (CO or H$_2$), mol m$^{-3}$
- $c_{i,TP}$: Concentration in liquid-filled transport pores of species $i$ (CO or H$_2$), mol m$^{-3}$
- $D_i$: Diffusion coefficient of species $i$ (CO or H$_2$), m$^2$ s$^{-1}$
- $E_{\text{Ab}}$: Activation energy for the reaction rate constant, J mol$^{-1}$
- $E_{\text{Ab}}$: Activation energy for the absorption constant, J mol$^{-1}$
- $H_i$: Henry coefficient of species $i$ (CO or H$_2$), bar
- $k_{ai}$: Combined mass transfer coefficient of species $i$ (CO or H$_2$), s$^{-1}$
- $k_s$: Selectivity coefficient, dimensionless
- $n$: Carbon number, dimensionless
- $P$: Gas phase pressure, bar
- $p_1$: Partial pressure of species $i$ (CO or H$_2$), bar
- $R$: Universal gas constant, J mol$^{-1}$ K$^{-1}$
- $r_{C_5+}$: Formation rate of desired products, mol s$^{-1}$ m$^{-3}$
- $r_{CO}$: Consumption rate of CO, mol m$^{-3}$ s$^{-1}$
- $S_{C_1}$: Spatially distributed molar selectivity for methane, dimensionless
- $S_{C_1,\text{eff}}$: Effective molar selectivity for methane, dimensionless
- $S_{C_5+}$: Spatially distributed molar selectivity for alkanes with $n$ greater than 4, dimensionless
- $S_{C_5+,\text{eff}}$: Effective molar selectivity for alkanes with $n$ greater than 4, dimensionless
- $S_{Cn}$: Spatially distributed molar selectivity for an alkane of carbon number $n$, dimensionless
- STY: Space time yield of desired products, mol s$^{-1}$ m$^{-3}$
- $t$: Thickness of catalyst layers, m
- $T$: Temperature, K
- $T_0$: Temperature at the wall boundary, K
- $T_{\text{max}}$: Maximum temperature in the layer, K
- $T_{\text{ref}}$: Reference temperature for selectivity expression, 493.15 K
- $v_l$: Molar volume of the products in liquid phase, m$^3$ mol$^{-1}$
- $x$: Dimension coordinate, m
- $H_2/CO$: Ratio of gas-phase partial pressures of H$_2$ to CO, dimensionless

**Greek**

- $\alpha$: Chain growth probability, dimensionless
- $\beta$: Selectivity order, dimensionless
- $\Delta E_a$: Selectivity expression activation energy, J mol$^{-1}$
- $\varepsilon_{\text{Cat}}$: Fraction of pores in the catalyst phase, dimensionless
- $\varepsilon_{TP}$: Fraction of transport pores in the catalyst layer, dimensionless
- $\eta_{\text{Cat}}$: Efficiency of the catalyst phase, dimensionless
- $\eta_{\text{Lay}}$: Efficiency of the catalyst layer including transport pores, dimensionless
- $\chi_{C_n}$: Stoichiometric coefficient of alkanes with $n$ greater than 29, dimensionless
- $\chi_{Cn}$: Stoichiometric coefficient of an alkane with carbon number $n$, dimensionless
- $\nu_i$: Stoichiometric coefficient of species $i$ (CO, H$_2$ or H$_2O$), dimensionless
- $\rho_{\text{Cat}}$: Apparent density of the catalyst phase, kg m$^{-3}$
- $\tau_{\text{Cat}}$: Tortuosity inside the catalyst phase, dimensionless
- $\tau_{TP}$: Tortuosity inside the transport pores, dimensionless

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