THz absorber for breast cancer early detection based on graphene as multi-layer structure

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Abstract
Today, the cost of treating cancer is very high at the medical center. So, early detection of the cancer is essential, and one of the possible ways to diagnose cancer in the early stages is to use sensors that are made in the terahertz area. The proposed absorber is a multilayer structure in which graphene is used in the upper and middle layer and the middle layer has a ribbon form. A disk form element is used for the upper layer. Finally, a ring is added to the structure. In the final absorber, the ring actually encompasses the field, and this will increase the Q-factor at 6.7 THz. The interaction of layers on each other and their effect on reflection is studied to modify the proposed absorber by variation of the chemical potential of the graphene layer in the range of 0.2 to 0.6 eV. Finally, the proposed absorber is used to discriminate cancer and healthy breast tissue. The effect of the thickness and distance of the tissue from the absorber is examined to realize the material effect on sensitivity and figure of merit (FOM) as two main factors for detecting the cancer tissue.

Keywords Absorber · Terahertz · Graphene · Cancer detection · Sensitivity

1 Introduction

Early cancer detection plays an important role for the medical centers in the microwave (Samsuzzaman et al. 2019), THz (Poorgholam-Khanjariand and Zarrabi 2021), and optical systems (Balaji and Zhang 2017). For this aim, various systems have been developed in the microwave spectrum for detecting breast cancer tumors during past decades (Shao and Adams 2013) and microwave antenna is the main element for these systems (Ruvio et al. 2013). However, by recent development in micro and nanofabrication, the THz and optical devices get more consideration for early cancer-detecting (Rahman et al. 2016). In fact, a small sample is sufficient for the THz spectroscopy system based on THz domain spectroscopy (TDS) (Yu et al. 2019) to recognize the cancer tumor from the normal tissue by analyzing either the pulse response (Woodward et al. 2002) or frequency shift (Kazemi 2021).
Moreover, the THz technology has been considered more for non-ionized characteristics (Keshavarz and Vafapour 2018) and therefore various devices have been designed such as antenna (Sadeghzadeh and Zarrabi 2016), filter (Azizi et al. 2018), metasurface (NouriNovin and Alomainy 2021), and absorber (Landy et al. 2008). Many models of the THz absorbers have been developed for different applications such as THz sensing (Fu et al. 2021). For example, split-ring resonators (Wen et al. 2009) and Jerusalem cross (Arezoomand et al. 2015) have been studied for THz application. The THz absorber is more interesting because the energy can be concentrated (Tavakoli et al. 2019). Thus the sensitivity and Figure of merit (FOM) will increase which are vital for the sensors (Soheilifar 2018).

On the other hand, various techniques based on special materials have been used to design the sensors (Mobasser et al. 2021), and using graphene is known as a conventional technique for developing THz absorber because of the plasmonic behavior of the graphene at this spectrum (Fu et al. 2018; Peng et al. 2019). Recently, many models of absorber by graphene have been investigated for narrowband (Cen et al. 2018; Li et al. 2016) or UWB spectrum (Biabanifard et al. 2019; Zhou et al. 2018).

The main purpose of this paper is to design an absorber for the terahertz spectrum at the 6.7 THz for breast cancer detection, which uses graphene in the design of this absorber and presents a multi-layered structure that has a high Q-factor capability. Then the normal and cancer tissue are examined for the proposed absorber. Finally, these cancerous and normal tissues are placed over the absorber structure and as a result, we obtain sensitivity and the Figure of merit for both tissues. So, three parameters of Q-factor, sensitivity, and FOM are determined for cancer and healthy tissue.

2 The graphene absorber designing

Graphene is a type of carbon allotropes. It is composed of a two-dimensional honeycomb crystal structure in which each carbon atom is attached to three other carbon atoms. The conductivity of graphene according to the Kubo formula has two parts, inter and intra, which are shown in the Eqs. 1 and 2, and the final conductivity of graphene can be obtained from the sum of these two parts (Zarrabi et al. 2016).

\[
\sigma_{\text{intra}}(\omega) = \frac{e^2K_B T}{\pi \hbar^2 (\omega - i2\Gamma)} \left[ \frac{\mu_c}{K_B T} + 2 \ln(e^{-\frac{\mu_c}{K_B T}} + 1) \right]
\]

(1)

\[
\sigma_{\text{inter}}(\omega) = \frac{e^2}{4\pi \hbar} \ln \left( \frac{2|\mu_c| - (\omega - i2\Gamma)\hbar}{2|\mu_c| + (\omega - i2\Gamma)\hbar} \right)
\]

(2)

Here, \(\sigma\) is the surface conductivity and \(e\) electric charge, \(\omega\) frequency angle, \(K_B\) Planck constant, \(T\) ambient temperature, \(K_B\) Boltzmann constant, \(\mu_c\) chemical potential, \(\tau\) relaxation time and \(\Gamma\) electron diffraction velocity and can be defined by \(\tau = 1/2\Gamma\).

The electrical conductivity of the graphene layer can be obtained by Eq. 3, where \(\sigma\) is the conductivity of the graphene layer, which is obtained from the Kubo equation of Eqs. 1 and 2, and \(\omega\) is the angular frequency and thickness \(d\) of the graphene layer (Soheilifar and Zarrabi 2019).
In Fig. 1a, we have compared the real part of surface impedance, which shows that for a chemical potential of 0.2 eV, the impedance is about 43 Ω, and this value has been reduced to 22 Ω for 0.4 eV. So it means that the surface shows lower impedance for higher chemical potential. As shown in Fig. 1b, increasing the relaxation time has a decreasing effect on the surface impedance of the structure like increasing the chemical potential of the graphene. Here, the chemical potential of the graphene is supposed 0.4 eV at 300°k for the thickness of 1 nm.

The proposed absorber is a multi-layered structure that at the bottom layer of a gold plate is located as a reflector. Then a quartz substrate is placed on the gold and graphene strips are placed on the first substrate. Then another layer of quartz material is placed on the graphene layer and finally, a graphene element is placed on it as the main element.

The proposed absorber is a multi-layered structure and the two absorbers are designed to be used as sensors. In fact, the difference between the two structures is in the main
absorber element. The proposed multi-layer structure has a gold layer in the bottom as a reflector with a thickness of 0.5 µm, which is used to eliminate the transmission, and on it is placed the main substrate with a thickness of 4 µm. The graphene ribbons with a width of 3 µm are placed on the first substrate and the distance between these ribbons is assumed to be 3 µm, and this layer is presented in Fig. 2a. Then we put a dielectric layer on the graphene ribbons and the thickness of this layer is assumed to be 3 µm and all the dielectrics are made of quartz with a permittivity of 3.75. Finally, a graphene element is placed on the second substrate. The boundary condition is one of the most important parts of an absorber simulation. Here, we have assumed perfect electrical conductor (PEC) in one direction and perfect magnetic conductor (PMC) in the other direction to provide periodic boundary conditions. For exciting the absorber, waveguide port excitation and periodic boundary conditions are used.

3 The simulation results and discussion

According to Fig. 3, the first structure has a resonance at the frequency of 6.7 THz. However, in the final structure, a disk and a loop have been used to improve the reflection, and the final structure has a resonance at same frequency with less reflection and is presented.
in Fig. 3. As can be seen, by using the ring structure, the amount of reflection is drastically reduced from $-15$ dB to $-20$ dB, which means having a stronger absorber as well as improving the Q-factor, which can be achieved through Eq. 4 (Ebrahimi and Poorgholam-Khanjari 2021) for the primary absorber with disk element the Q-factor is obtained 419.4 and for the proposed absorber when the ring added the Q-factor increases to 670.

$$Q = \frac{f_0}{BW_{3dB}}$$

Two factors, chemical potential and relaxation time will affect the frequency response of our structure. First, we consider the chemical potential of the ribbons as 0.2 eV and the changes in the chemical potential of the main element (the disk and the disk and ring) for the chemical potentials of 0.2, 0.4, and 0.6 eV. As shown in Fig. 4a for a chemical potential of 0.2 eV of the disk for the primary absorber, the structure has two resonances, but

**Fig. 4** The reflection of the absorber with chemical potential of 0.2 eV for graphene ribbon and various chemical potential of the main element a for primary absorber, b for final absorber
the reflection is not very significant and the reflection is about $-9$ to $-11$ dB. Then, with increasing chemical potential to 0.4 and 0.6 eV, the primary absorber has one resonance, but it has a lower reflection which reaches about $-15$ to $-17$ dB. Figure 4b shows that the final structure has dual-band properties for a chemical potential of 0.2 eV for the disk and ring with a chemical potential of the 0.2 eV for ribbons, but both resonances are in the frequency range of 7 to 8 THz, but by increasing chemical potential, the structure has a resonance at lower frequencies in it happens about 6.7 THz. As shown in Fig. 1a, the surface impedance of the graphene layer will change by variation of the chemical potential and this change can be considered for controlling the resonances. Of course, the interaction of the top layer and middle layer plays important role in the proposed structure.

To understand the effect of graphene ribbons, we assumed the disk potential to be 0.6 eV and changed the chemical potential of the graphene ribbons $e$ to 0.2, 0.4, and 0.6 eV, and the results are shown in Fig. 5a. Here it is observed that with the increasing chemical potential of the graphene ribbons, the resonances are transferred to a higher frequency of 7.5 THz, and the reflection is increased to $-10$ dB, so the chemical potential of

![Fig. 5](image.png)

The reflection of the absorber with chemical potential of 0.6 eV for the main element and various chemical potential of the graphene ribbon a for primary absorber, b for final absorber
the graphene ribbons is reduced to 0.2 eV gives the best reflection. When we assume that the chemical potential of the ring and the disk is constant and the chemical potential of the graphene ribbons changes, we observe behavior in Fig. 5b that we saw earlier in Fig. 5a for the structure of a disk without a ring, except that the ring reduces the reflection from the absorber. In fact, the loop prevents the scattering of fields and thus creates a capacitor around the disk, causing energy storage as a result of increasing the Q factor of the structure.

As shown in Fig. 6, we examine the effect of relaxation time on the reflection of the absorber. We checked the relaxation time of graphene for the disk and ring in the final structure for the values of 0.5, 1 and 1.5 pS. As shown in the results, the relaxation time changes have no effect on the operating frequency, while the reflection value changes drastically, which also affects the Q-factor. The Q-factor of the proposed absorber for the relaxation time of 0.5 pS is 239 but this amount is 670 and 1970 for relaxation time of 1 pS and 1.5 pS respectively. So, increasing the relaxation time can improve the Q-factor of the absorber. As shown in Fig. 1b, the surface impedance of the graphene layer will change by variation of the relaxation time and this change is noticed to modify the return loss (reflection) from the proposed surface. In fact, the impedance of the surface will reduce by increasing the relation time and therefore, the reflection from the surface will decrease.

In Fig. 6, the distribution of electric and magnetic fields for the structure of the disk at a frequency of 6.15 THz at the absorber surface and the interface of the graphene ribbons and substrate is investigated. Figure 7a shows the value of electric field near graphene ribbons, which is about 136 dB, which is the maximum field in the X direction, and Fig. 7b shows the value of electric field at the absorber surface. It shows that it is about 148 dB, and here the maximum field is clearly distributed in the X direction with dipole form. Figure 7c shows the magnitude of the magnetic field near graphene ribbons, which is about 90.7 dB. In Fig. 7d shows the amount of magnetic field at the absorber surface, which is about 95 dB, and here the maximum field is clearly distributed inside the disk and actually creates an inductance property.

In Fig. 8, the distribution of electric and magnetic fields for the ring structure (final absorber) at a frequency of 6.70 THz at the absorber surface and the interface of the graphene ribbons and substrate is investigated. Figure 8a shows the electric field on the graphene ribbons, which is about 122 dB. Figure 8b shows the electric field at the absorber surface.
surface, which is about 139 dB, and here the maximum field is clearly distributed in the X direction and dipole form. Figure 8c shows the magnitude of the magnetic field on the graphene ribbons, which is about 87.8 dB and Fig. 8d shows the magnitude of the magnetic field at the absorber surface, which is about 85 dB.

### 4 Cancer detection with THz absorber

As shown in Eq. 5, the Debye second order can be used to model breast and skin cancer and tissue (Vafapour et al. 2020):

\[
\varepsilon(\omega) = \varepsilon_\infty + \frac{\varepsilon_S - \varepsilon_2}{1 + i\omega\tau_1} + \frac{\varepsilon_2 - \varepsilon_\infty}{1 + i\omega\tau_2} \tag{5}
\]

The \(\varepsilon_\infty\), \(\varepsilon_S\), \(\varepsilon_2\), \(\tau_1\) and \(\tau_2\) are temperature dependent of high frequency permittivity, static dielectric constant, intermediate frequency limit, slow relaxation time, and fast relaxation time. The parameters for breast cancer and breast healthy tissue are given in Table 1.

In Fig. 9, the variation of the real and imaginary parts of the permittivity for healthy and cancerous breast tissue based on the second-order Debye model in the range of 4 to 8 THz. As can be seen, the cancerous tissue has a higher permittivity in the real and imaginary part, so we expect more frequency shifts and losses in the results for the cancerous tissue.
Here, we have considered two factors for cancer diagnosis and compared this issue for healthy and cancerous tissue. For breast cancer tissue, the results are presented in Fig. 10. Here we assume the distance of the object from the absorber is 0, 5 and 10 µm, respectively. As can be seen in the figure, the frequency shift will be more by increasing distance for a constant thickness of 1 µm of the sample tissue. When the tissue is placed on the absorber surface, the operating frequency is about 6.4 THz, but with increasing height, the working frequency has reached 5.8 THz. The reason for this phenomenon is the appearing a capacitance between the sample tissue and the absorber, which reduces the operating frequency. The sensitivity and FOM of the sensor can be used for determining the healthy and cancer cell when the tissues are placed in front of them. The sensitivity can be obtained by Eq. 6 where the Δf is the frequency shift and the variation of the refractive index is indicated by Δn (Vafapour et al. 2020; Poorgholam-Khanjari et al. 2021).

| Table 1 | Debye coefficients for healthy and cancerous breast tissue |
|---------|----------------------------------------------------------|
|         | ε∞  | ε₅     | ε₂   | τ₁[Ps] | τ₂[Ps] |
| Healthy tissue | 2.1  | 76.5   | 3.9  | 10.3   | 0.07   |
| Cancer tissue   | 2.5  | 77.9   | 4.3  | 9.1    | 0.08   |

Fig. 8 Field distribution for proposed structure a electric field on the graphene ribbons, b electric field in the absorber surface, c magnetic field on the graphene ribbons, d magnetic field in the absorber surface
Fig. 9 The real and imaginary parts of the permittivity for healthy and cancerous breast tissue

![Graph showing the real and imaginary parts of the permittivity for healthy and cancerous breast tissue.]

Fig. 10 Effect of tissue distance from absorber a healthy tissue for breast tissue, b cancer tissue for breast tissue

![Graph showing the effect of tissue distance on return loss for (a) healthy tissue and (b) cancer tissue.]
\[ S = \frac{\Delta f}{\Delta n} \]  \hspace{1cm} (6)

And the FOM can be calculated by Eq. (7), where the \( S \) is the sensitivity of the sensor and FWHM (Full width at half maximum) is the 3 dB bandwidth of the reflection (Vafapour et al. 2020; Poorgholam-Khanjari et al. 2021).

\[ FOM = \frac{S}{FWHM} \]  \hspace{1cm} (7)

In Fig. 11, the effect of tissue thickness change on the frequency response is examined. The healthy and cancerous breast tissues with thicknesses of 1, 3 and 5 \( \mu m \) are placed on the absorber surface. Obviously, by increasing thickness, the capacitance will increase and as a result, the frequency shift will increase. Also, the frequency shift is higher for cancerous tissue than for healthy tissue.

![Fig. 11 Effect of tissue thickness on absorber a healthy tissue for breast tissue, b cancer tissue for breast tissue](image-url)
The result for sensitivity and FOM for each case of our studied based on Eqs. 6 and 7 are obtained and the presented in Table 2 for various distance. Then, the effect of thickness on sensitivity and FOM are presented in Table 3. As shown in Tables 2 and 3, the FOM and sensitivity reveal the discrimination between the healthy and cancer tissue clearly and we can use it for detecting the cancer very fast and accurate in the THz spectrum with a small sample. The sensitivity of the sensor will increase by increasing the thickness and distance of the tissues as shown in Tables 2 and 3.

Typically, the gold (Arezoomand et al. 2015) or graphene element are using for the THz absorber and the THz absorber can be put in two main categories of wide (Biabanifard et al. 2019; Zhou et al. 2018) and narrow bandwidth (Cen et al. 2018; Li et al. 2016). The main benefit of the proposed absorber in comparison with gold base THz absorber is the reconfigurable characteristic of this absorber while and we can control the Q-factor with relaxation time while this option is not possible in conventional THz absorber (Arezoomand et al. 2015; Fu et al. 2021). Moreover, some studies about THz absorber with graphene have been focus on wide bandwidth absorber while for sensing, higher Q-factor is benefit for detecting material with higher FOM (Biabanifard et al. 2019; Zhou et al. 2018).

5 Conclusion

A narrowband THz absorber was developed for cancer breast detection and the operation frequency of the proposed absorber is 6.5 THz. As the first point, the effect of the layer to improve the Q-factor of the absorber was shown and the ring element plays important rule for trapping the electric field in a limited area for the proposed absorber. The Q-factor of the absorber was improved from 419 for primary model to 670 in the proposed absorber. The results show the effect of thickness and distance of the tissue sample from absorber on sensitivity and FOM which can be considered to discriminate the cancer and healthy tissue. In fact the proposed absorber shows high FOM which is about 29,117 RIU$^{-1}$ whereas the FOM is 10344 for the healthy tissue.

| Table 2 | The sensitivity and FOM of the sensor for healthy and cancer tissue for various distance |
|---------|-----------------------------------------------|
|         | L = 0 μm | L = 5 μm | L = 10 μm |
| Sensitivity (GHz/RIU) | Healthy tissue | 82 | 289 | 379 |
|               | Cancer tissue | 101 | 292 | 458 |
| FOM (RIU$^{-1}$) | Healthy tissue | 10,344 | 24,137 | 9482 |
|               | Cancer tissue | 29,117 | 18,312 | 7643 |

| Table 3 | The sensitivity and FOM of the sensor for healthy and cancer tissue for various thickness |
|---------|-----------------------------------------------|
|         | d = 1 μm | d = 3 μm | d = 5 μm |
| Sensitivity (GHz/RIU) | Healthy tissue | 82 | 296 | 475 |
|               | Cancer tissue | 10,344 | 9267 | 5172 |
| FOM (RIU$^{-1}$) | Healthy tissue | 101 | 335 | 590 |
|               | Cancer tissue | 29,117 | 10,470 | 1605 |
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