Stimulated electromagnetic emission polarization under different polarizations of pump waves

E. D. Tereshchenko¹, R. Y. Yurik¹, and L. Baddeley²,³

¹Polar Geophysical Institute, Murmansk, Russian Federation
²University Centre on Svalbard, Longyearbyen, Norway
³Birkeland Centre for Space Science, University of Bergen, Bergen, Norway

Correspondence to: R. Y. Yurik (roman.yurik@pgi.ru)

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Abstract. The results of investigations into the stimulated electromagnetic emission (SEE) polarization under different modes of the pump wave polarization are presented. The present results were obtained in November 2012 during a heating campaign utilizing the SPEAR (Space Plasma Exploration by Active Radar) heating facility, transmitting in both O- and X-mode polarization, and a PGI (Polar Geophysical Institute) radio interferometer capable of recording the polarization of the received radiation. The polarization ellipse parameters of the SEE DM (downshifted maximum) components were determined under both O-mode and X-mode polarization of the pump waves. The polarization direction of the SEE DM component was preserved under different polarizations of the pump waves. Different polarizations of the pump waves have a different SEE generation efficiency. The intensity of the DM component is observed to be greater during O-mode pumping. In addition, the numbers of observed SEE features are also greater during O-mode pumping.

Keywords. Ionosphere (active experiments; plasma waves and instabilities) – radio science (ionospheric physics)

1 Introduction

Since the first registration of stimulated electromagnetic emission (SEE), during the heating experiments at Tromsø (Thiedé et al., 1982), this phenomenon has been investigated using mainly amplitude methods. The SEE is a weak signal, limited across a frequency band of ~150 kHz around the pump frequency which is generated through the excitation of the ionospheric plasma by powerful radio emissions from a high-frequency (HF) heating facility. Observations of the SEE under different ionospheric conditions and different operation modes of the pump wave allow the characteristic features of the spectrum to be distinguished and the conditions affecting their generation to be identified.

It is known that the artificial generation of spectral components depends on several factors, such as the polarization of the pump wave, ionospheric conditions and the proximity of the pump wave to the local harmonic of the cyclotron frequency (gyroharmonic) of electrons (Stubbe et al., 1984; Leyser, 2001). The spectral width of the various components of the SEE may be a few hertz in the case of narrow-band components near the pump wave frequency and greater than 100 kHz in the case of broadband radiation components.

An investigation into the SEE using radio-interferometric methods was started in 2000. This study has indicated that the processes involved in SEE generation are more complex than previously thought. In some cases the region of SEE generation may not coincide with the direction of maximum radiation of the pump wave (Carozzi et al., 2001; Tereshchenko et al., 2006). Additionally, Tereshchenko et al. (2012) showed that the emission of a discrete spectral feature of the SEE is not from a single-point location.

The study of the polarization characteristics of SEE near the fourth harmonic of the electron cyclotron frequency was carried out using the heating facility “Sura” in 1998. Studies have shown (Carozzi et al., 2001) that this type of measurement can be a new and informative diagnostic tool for SEE. In particular, it was found that part of the SEE spectrum has a low degree of circular polarization, which may be caused by the presence in the observed signal of overlapping different components with left and right polarizations and by the fact...
that the SEE signal comes from the spatial distribution of the
generation region. Unfortunately, due to the sporadic nature
of the study of data, many questions about the polarization of
the SEE were not resolved in the course of the experiment.

The X-mode pumping has been performed in a number of
experiments, with and without SEE diagnostics. Even in
cases when no SEE diagnostics were used, the results bear
on the physics of the ongoing processes in the present ex-
periments (e.g. Thidé et al., 1983; Robinson et al., 1997). In
this paper we present the results from a heating campaign
utilizing the SPEAR (Space Plasma Exploration by Active
Radar, Robinson et al., 2006) heating facility (located near
Longyearbyen, Svalbard), which transmits circular polariza-
tion, and a radio interferometer capable of recording the
polarization of the received radiation (located in Barentsburg,
Svalbard).

Combined interferometric and polarization measurements
allow us to investigate not only the general characteristics of
the polarization of the received radiation but also to deter-
mine the radiation characteristics from different arrival di-
rections.

2 Description of the experiment

The SPEAR heating facility is situated on Spitsbergen in
the Svalbard archipelago (78.15° N, 16.06° E). SPEAR is de-
signed to carry out a range of space plasma investigations
into the polar ionosphere and magnetosphere. The SPEAR
site is located adjacent to the EISCAT (European Incoher-
ent Scatter Scientific Association) Svalbard Radar (ESR).
The SPEAR antenna system comprises a 6 × 4 array of full-
wave, crossed dipoles, 16 m above the ground, with an an-
tenna spacing of 48.4 m, allowing the transmission of both
linear and circularly polarized signals. The individual dipoles
are rhombically broadened to allow operation between 4 and
6 MHz. The resulting beam has a quasi-elliptical cross sec-
tion, with an average half-power width of 21° along its ma-
ajor axis and 14° along its minor axis. This results in an over-
all antenna gain of 21 dB. Individual phase control provides
beam steering within ±30° from zenith, at any azimuth, with-
out significantly altering the antenna gain or introducing sig-
ificant side-lobe signals. The beam may thus be pointed in
directions which include vertical and field-aligned direc-
tions. SPEAR has a distributed transmitting and receiving
system consisting of individual 4 kW solid-state transmis-
ters connected to the antenna array. Each transmitter con-
ists of a single driver stage and four individual 1 kW mod-
ules, the outputs of which are combined and fed to the an-
tenna. In the current configuration 48 such transmitters are
required. These are made up of 192 power-amplifier mod-
ules, 48 driver stages and a total of 240 power-supply units.
Each transmitter contains an embedded controller and an in-
dividual direct digital synthesizer (DDS) which provides a
transmit signal and the IF for the receiver front end.

For the polarization measurements of the SEE presented
here, two channels of an HF interferometer designed at the
Polar Geophysical Institute (PGI) were employed. This in-
terferometer has previously been used for the determination
of the angle-of-arrival of SEE signals at the Sura (Nizhniy
Novgorod, Russia), EISCAT (Tromsø, Norway) and SPEAR
heating facilities (Isham et al., 2005; Tereshchenko et al.,
2006, 2008, 2011). The HF interferometer consists of an
HF receiver with a high dynamic range and a bandwidth
of 300 kHz, a digital system for the conversion and record-
ing of the received signal, and HF and IF (intermediate fre-
quency) notch filters. The two channels of the HF receiver
were used for the polarization measurements. The receiv-
ing antenna consists of a dipole antenna with 9 m masts which
are oriented in the north–south and west–east geomagnetic
planes. The data logging system uses two 16 bit analogue-
digital converters (ADC) with two channels on each and a
sampling frequency of 60 MHz. All channels of the ADC
can work in synchronous and nonsynchronous modes. The
HF and IF notch filters can be used for enlarging the dy-
namic range of the ADC. The receiving point was located
at a PGI observatory (coordinates: 78.094° N, 14.208° E) lo-
cated about 40 km south-west of the SPEAR site.

3 The calculation and data processing methods

The coordinate system used for the Northern Hemisphere is
shown in Fig. 1. On the figure the x axis is directed toward
the north, and the y axis is directed toward the east. When
using a dual-channel receiver for an orthogonal arrangement
of dipole antennas oriented in the geomagnetic meridional
(north–south) and zonal (east–west) planes (Fig. 1), the re-
ceived signal of the quasi-monochromatic wave can be rep-
resented as

\[
\begin{align*}
E_x &= E_{x0} \cos(\omega t + \Delta \varphi) \\
E_y &= E_{y0} \cos(\omega t + \Delta \varphi + \Delta \varphi) 
\end{align*}
\]

where \(E_{x0}\) and \(E_{y0}\) are the amplitude components of
the electric oscillation and \(\Delta \varphi\) is the phase difference
between the oscillations in the meridional and zonal direc-
tions.

In those cases where \(\Delta \varphi = n \pi (n = 0, 1, 2, \ldots)\), the ob-
served oscillations are linearly polarized. In the case of
circular polarization, the phase difference \(\Delta \varphi = (2m - 1)\pi/2(m = 1, 2, 3, \ldots)\).

The electric field vector in the horizontal plane describes
an ellipse which can be represented by the equation

\[
x'^2/a^2 + y'^2/b^2 = 1,
\]

where \(x'\) and \(y'\) represent coordinates in an orthogonal coor-
dinate system oriented along the ellipse axes; \(a\) and \(b\) are
the minor and major semi-axes of the ellipse:

\[
a^2 = \frac{D^2}{2A + 2\sqrt{A^2 - D^2}},
\]

\[
b^2 = \frac{D^2}{2A - 2\sqrt{A^2 - D^2}}.
\]
The phase difference between the electric field components in the zonal and meridional planes ($\Delta \varphi$) was determined from the ratio between the real and imaginary parts of the complex parameter $E_{xy}$, i.e.

$$\Delta \varphi = \arctg \left( \frac{D}{C} \right) = \arctg \left( \frac{\text{Im}[E_{xy}]}{\text{Re}[E_{xy}]} \right).$$

where $E_{xy}$ represents the complex parameter of the two channels, $E_x$ and $E_y$.

The calculated absolute value of squared coherence (SC) of a two-channel system was used to differentiate between the radiated polarized components:

$$SC = \left| \frac{E_{xy}}{\sqrt{E_x E_y}} \right|^2.$$  \hspace{1cm} (10)

In a real signal, where there is a noncorrelated noise, $SC \neq 1$.

4 The results

This paper focuses on the observation of the SEE polarization parameters under different polarizations of pump waves. While the ADC utilized a sampling frequency of 60 MHz, each spectrum was obtained using an integration time of 2.2 s during one record. Twenty-seven series of observations (each series contains 70 records for 11 min of observation) which were obtained during the heating campaign from 13 to 17 November, were used for investigations. We selected four series which correspond time periods of similar ionospheric conditions and SPEAR operation mode. However, we have selected only that registration for which we observed a good SEE signal. The data were selected from four experiments carried out during 2012: two using O-mode polarization of the pump wave (13 November, 10:39:56 UTC, and 16 November, 09:33:26 UTC) and two using X-mode polarization (13 November, 11:25:26 UTC, and 16 November, 10:54:11 UTC).

The SPEAR radiated the pump waves along the geomagnetic field line (into magnetic zenith with an azimuth of 184° and with an antenna main lobe inclination of 8° south). The effective radiated power (ERP) was about 17 MW. The pumping series includes a 5 min pause before radiation was switched on and 10 min of radiation.

The possible mixture of both O and X polarizations in the transmission is a very important factor which should be taken into account for experiments and further interpretation of the polarization measurements. The most critical factors here are the phasing of each pair of crossed dipoles and the matching of the two antennas in power. This imbalance will give some leakage in the opposite mode, but it is a very small value and can be a few tens of decibels less than the main radiation. However, it was found early on that threshold for the SEE generation is about 5 MW ERP (Leyser et al., 1994; Frolov et al., 2000). Taking into account the low power of the SPEAR
radiation, we have assumed the absence of another polarization which was capable of exceeding the threshold of SEE generation.

It should be noted that SEE has been detected with less than 100 kW ERP of pump wave during preconditioning from previous pumping (e.g. Leyser et al., 1990). However, for the SPEAR location the horizontal plasma drift can remove the pump-induced irregularities from the pump–ionosphere interaction region (Tereshchenko et al., 2011). Figure 2 shows the PSD observed during the experiments conducted on the 13 (upper panel) and 16 (lower panel) November 2012 using the O-(left panel) and X-mode polarization (right panel) of the pump wave. The pump frequency is strongly attenuated by the notch filters and is not shown in the figures.

In all four cases, the DM (downshifted maximum) feature of the SEE spectrum was observed. The intensity of the DM was 20 dB above the noise floor in the case of O-mode polarization and about 10 dB in the case of the X-mode polarization of the pump wave.

When using the pump wave with O-mode polarization, the BC (broad continuum) component is observed in the generated SEE spectrum. Spectral widths of the BC were up to 50 kHz on 13 November and 40 kHz on 16 November. During O-mode polarization the DM intensity is comparable between the 13 and 16 November; however, the BC intensity on the 16 November is decreased by \( \sim 5 \) dB.

Additionally, the SEE spectra observed during O-mode polarization of the pump wave on the 13 November showed an upshifted maximum (UM) feature at a frequency of \( \Delta f = 8.6 \) kHz above the pump wave frequency (4.45 MHz) with an intensity of about 5 dB above the noise level.

When the X-mode polarization of the pump wave was used, the BC was not observed. The UM feature at a frequency of 9.3 kHz with an intensity exceeding the floor noise level by 3 dB was observed on 16 November. In both cases, when a pump wave with X-mode polarization was used, the observed intensity of SEE was much smaller in magnitude when compared to the cases when a pump wave with O-mode polarization was used.

It should be noted that the polarization of the propagating waves is preserved only when the plasma density gradient is parallel or antiparallel to the geomagnetic field, which generally is not the case in the ionosphere. Thus, in general, all transmitted circularly polarized waves will be decomposed into two oppositely polarized waves (of different amplitudes) when entering into the ionosphere. In cases when wave propagation is not along geomagnetic field lines in the SPEAR beam, the transmitted circularly polarized wave will decompose into elliptically polarized O-mode and X-mode waves when entering into the ionosphere. For cases when SPEAR transmits pump waves along the geomagnetic field lines (slope angle equals \( 8^\circ \) and azimuth equals 184\(^\circ\)), the angle between transmitted ray and geomagnetic field line can achieve 9\(^\circ\). Thus, at the edges, the pump wave pattern has some ellipticity.

For further analysis the spectral bandwidth from the pump wave frequency to 40 kHz below the pump wave frequency was used (marked by the vertical dashed line in the plots in Fig. 2). Figure 3 shows plots of PSD (top row) and the coherence of the polarized component (SC) (bottom row) in the frequency range from \(-20\) to 0 kHz relative to the pump frequency. The PSD includes the DM feature. The analysis of the polarization parameters was carried out for the DM component because this spectral component of the SEE had been observed in all cases.
Plots of the squared coherence (Fig. 3, bottom panels) show that, in all cases, the DM component of the SEE has a high degree of coherence and can be distinctly distinguished in the received signal except on 13 November 2012, when BC had a similarly high coherence. The absence of the highest values of coherence at 10:39:56 UT on 13 November 2012 (bottom row, first plot) could be caused by the more extended excitation region, which in turn gives less coherence because of the superposition of emissions at the receiving site. In the case of the pump wave with O-mode polarization, the observed BC feature of the SEE also has a high coherence.

The DM component of the SEE is located between $-5$ and $-15$ kHz from the pump wave frequency. In the case of O-mode polarization, a squared coherence level of 0.65 was observed. In the case of the X-mode polarization, a squared coherence level of 0.45 and 0.35 was observed for the experiments conducted on the 13 and 16 November respectively.

Figure 4 shows the phase difference $\Delta \varphi$, the angle of rotation of the polarization ellipse $\theta$ (upper panel) and eccentricity $e$ (lower panel) of the DM component of the SEE for all the selected four cases.

In all cases, the calculated eccentricity of the ellipse of polarization lies in the range of 0.8 to 0.9, indicating an elliptical polarization of the observed signals, which arrive from the disturbed region according to the interferometer data. The angle between the major axis of the polarization ellipse and the meridional plane lies in the range from $-10$ to $20^\circ$, i.e. the polarization ellipse is focused mainly in the geomagnetic meridian plane. The slight differences in the values of the angle between the major axis of the polarization ellipse and the meridional plane indicate possible differences in the ionosphere or magnetic conditions at the time of observation, which lead to changes in the location of the pump–plasma interaction region and the conditions of the wave propagation from a one series of observations to another and from day to day.

The mean values of $\Delta \varphi$ are close to $90^\circ$ and were $100$ and $87^\circ$ during O-mode polarization (10:39:56 UTC on 13 November and 09:33:26 UTC on 16 November respectively) and 79 and $118^\circ$ during X-mode polarization (11:25:26 UTC on 13 November and 10:54:11 UTC on 16 November respectively).

The variation of $\Delta \varphi$ observed during periods of X-mode polarization is greater than that observed during O-mode polarization. It could be related to the weaker signal of the SEE under the X-mode pumping. Also, the variation of parameters may be due to differences in the ionospheric conditions during the observations as a result of high variability of ionospheric parameters in the polar cap.

5 Discussion

The utilization of the different pump wave polarizations leads only to a change in the efficiency of the generation of the SEE that was found during the early phase of SEE investigations (Thidé et al., 1983). The predominant O polarization of DM was preserved in both cases of the pump wave polarizations used. The preservation of the polarization direction of DM under different polarizations of the pump waves may indicate that the DM signals are generated in the same way.

The observed polarization ellipses of the DM signals were oriented in such a way that their major axes were in the meridional plane while, according to the geometry of the experiment, the arrival direction of the SEE signals was east. It is possible that the ellipticity of the polarization of the SEE was associated with the geometry of the signal propagation but was also affected by the spatial elongation of the generation region along the magnetic field. Thus, as the results show, in the magnetized ionospheric plasma the observed radiation of DM is elliptically polarized.

Although the theoretical concept of polarization was formulated mainly for quasi-monochromatic plane waves, the results show that this consideration is applicable to the study of the polarization characteristics of the particular spectral components (at least narrow maxima and peaks).

It is suggested that further studies of the polarization characteristics of the SEE should be carried out in combination with interferometric and polarization measurements in addition to other instrumentation utilized in the study of plasma dynamics. This would allow the determination of the radiation characteristics from different spatially localized sources.

6 Conclusions

Measurements of the polarization characteristics of SEE under different pump wave polarizations were undertaken and
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