An artificial neural network framework for reduced order modeling of transient flows

Omer San*, Romit Maulik, Mansoor Ahmed
School of Mechanical and Aerospace Engineering, Oklahoma State University, Stillwater, Oklahoma 74078, USA

Abstract
This paper proposes a supervised machine learning framework for the non-intrusive model order reduction of unsteady fluid flows to provide accurate predictions of non-stationary state variables when the control parameter values vary. Our approach utilizes a training process from full-order scale direct numerical simulation data projected on proper orthogonal decomposition (POD) modes to achieve an artificial neural network (ANN) model with reduced memory requirements. This data-driven ANN framework allows for a nonlinear time evolution of the modal coefficients without performing a Galerkin projection. Our POD-ANN framework can thus be considered an equation-free approach for latent space dynamics evolution of nonlinear transient systems and can be applied to a wide range of physical and engineering applications. Within this framework we introduce two architectures, namely sequential network (SN) and residual network (RN), to train the trajectory of modal coefficients. We perform a systematic analysis of the performance of the proposed reduced order modeling approaches on prediction of a nonlinear wave-propagation problem governed by the viscous Burgers equation, a simplified prototype setting for transient flows. We find that the POD-ANN-RN yields stable and accurate results for test problems assessed both within inside and outside of the database range and performs significantly better than the standard intrusive Galerkin projection model. Our results show that the proposed framework provides a non-intrusive alternative to the evolution of transient physics in a POD basis spanned space, and can be used as a robust predictive model order reduction tool for nonlinear dynamical systems.

Keywords: Artificial neural networks; reduced order modeling; proper orthogonal decomposition; convective flows; non-intrusive model order reduction

1. Introduction
The full resolution of the dynamics of engineering flows requires a computational expense that is as yet infeasible for any practical purpose. Due to these overwhelming computational

*Corresponding author.
Email address: osan@okstate.edu (Omer San)
demands, model order reduction type approaches are gaining popularity. They are extensively used for many optimization and control applications (Akhtar et al., 2015; Buffoni et al., 2006; Dyke et al., 1996; Fang et al., 2009; Fortuna et al., 2012; Freund, 1999; Noack et al., 2011; Lucia et al., 2004; Roychowdhury, 1999; Silveira et al., 1996). Among the multitude of reduced order modeling approaches, proper orthogonal decomposition (POD) has emerged as a popular technique for the study of dynamical systems (Aubry et al., 1988; Holmes et al., 1998; Ly and Tran, 2001; Kerschen et al., 2005; Cizmas et al., 2008; Amsallem and Farhat, 2012; Taira et al., 2017). To clarify, POD is also documented under terminology such as the Karhunen-Loève expansion (Loève, 1955), principal component analysis (Hotelling, 1933) or empirical orthogonal function (Lorenz, 1956). POD extracts the most energetic modes from a collection of high fidelity numerical simulations of the governing equations of the dynamic system being studied. These bases are then used to reduce the degrees of freedom of the governing equation to scales that are computationally tractable. In particular, the fluid mechanics community has traditionally used POD as a method of extraction of the large scale coherent structures (represented by these modes) for the purpose of statistical pattern recognition (Lumley, 1967). This is because the global POD modes optimally span our physical space through a considerably truncated number of bases to resolve attractors and transients well.

The evolution equations for the lower order system are then obtained using the Galerkin projection (GP) method. The POD-GP method has been extensively used to provide fast and accurate simulations of large nonlinear systems (Berkooz et al., 1993; Kunisch and Volkwein, 2002; Lucia et al., 2004; Borggaard et al., 2007; Weller et al., 2010; Benner et al., 2015; Brunton and Noack, 2015; Taira et al., 2017; Rowley and Dawson, 2017). The Galerkin projection is devised so as to generate a reduced model where the truncated linear coefficients of the POD bases are evolved through time as a set of ordinary differential equations (ODEs). This process is denoted a projection since it is obtained by an orthogonal projection of the governing equations onto the truncated POD modes. The result is a considerably truncated system which continues to represent the most significant characteristics of the governing equations. The small memory requirement and low computational expense of this framework make it particularly suitable for post-processing analysis, optimization and control type problems where repeated model evaluations are required over a large range of parameters. This approach is significantly effective for quasi-stationary, time-periodic and decaying problems but might be challenging for highly non-stationary, convective and nonlinear problems.

The drawbacks are generally due to the lack of adaptation of the POD modes as well as the loss of information due to the truncation being limited to the most energetic of modes (Cordier et al., 2013; El Majd and Cordier, 2016). Amongst the various efforts used to mitigate these drawbacks include modeling the effect of truncated modes on the retained ones, an approach known as closure modeling (Wang et al., 2011; Borggaard et al., 2011; Wang et al., 2012; Östh et al., 2014; San and Iliescu, 2015; Wells et al., 2017; Xie et al., 2017; 2018b, a), and developing strategies to construct a more representative basis (Fillo et al., 2000; Buffoni et al., 2006; Bui-Thanh et al., 2007; Kunisch and Volkwein, 2010; Carlberg and Farhat, 2011; San and Borggaard, 2015; Ahmed and San, 2018). On the other
hand, the energy balance and mass conservation properties of projection based reduced order models have also been studied for fluid flows (Mohebujjaman et al., 2017). However, the projection-based model reduction approaches have limitations especially for complex systems such as general circulation models, since there is a lack of access to the full-order model (ROM) operators or the complexity of the forward simulation codes that renders the need for obtaining the full-order operators (Peherstorfer and Willcox, 2016). Therefore, there is a recent interest in generating fully non-intrusive approaches without the need for access to full-order model operators to establish surrogate models (Audouze et al., 2013, Mignolet et al., 2013, Xiao et al., 2015b, Hesthaven and Ubbiali, 2018, Hampton et al., 2018, Chen et al., 2018b, Xiao et al., 2019, Wang et al., 2019).

Alternatively, machine learning offers promise to generate accurate parametric reduced order models. There are a broad range of opportunities for the development of new deep learning architectures for reduced order models and numerical analysis of these machine learning approaches in applications of challenging flows. Recently, a supervised machine learning approach has been used to develop feasible regions in the parameter space where the admissible target accuracy is achieved with a predefined reduced order basis (Moosavi et al., 2015). Neural networks based closure schemes have also been developed by the authors to compute a stabilization term (i.e., eddy viscosity) for Galerkin projection based reduced order models (San and Maulik, 2018a, b, c). A deep residual recurrent neural network has been introduced as an efficient model reduction technique for nonlinear dynamical systems (Kani and Elsheikh, 2017). In a context similar to this recurrent neural network architecture, we formulate our problem in low-dimensional embedded structures that capture the major portion of the energy of the system in order to take advantage of lower-dimensional coherent structures that are responsible for the bulk mass, momentum, and energy transfer instead of performing a very large-scale simulation that considers all structures. A recent investigation of a spectral POD (SPOD) spanned model order reduction framework for fluid flows has been performed by Lui and Wolf (2019) using the deep feedforward neural networks. A long short-term memory (LSTM) architecture based model reduction approach has also been recently introduced to complement an imperfect reduced order model with data-streams utilizing a regularized hybrid framework (Wan et al., 2018).

The present study aims to bypass the Galerkin projection through the use of a single-layer artificial neural network (ANN) by introducing two different architectures, namely a sequential network (SN) and a residual network (RN). ANNs can be classified under the broader category of machine learning methods (i.e., systems that learn from data) and are a mathematical representation of the biological neural networks found in the human central nervous system. They have been used for a wide variety of applications such as function approximation, classification, data processing and dynamic systems control (Widrow et al., 1994). Briefly, an ANN can be used to setup a nonlinear relationship between a desired set of inputs and targets provided a large amount of benchmark data for the underlying relationship is available. This fitting to available data (also known as training) ideally generates a set of linear combinations of parameters and transfer functions which replicate the mean behavior of the underlying phenomena. This allows for the representation of subtle relationships which cannot be expressed explicitly in a functional form. The interested reader is directed
to the excellent introductory text by Demuth et al. (2014) on the development of a variety of ANN architectures and their underlying principles. We stress here that generating a robust non-intrusive reduced order modeling (ROM) approach is highly desirable for its ability to preclude model-form uncertainty in our forward simulations for reduced computational degree of freedom and is thus a highly active area of research (Xiao et al., 2015; Bistrian and Navon, 2015; Peherstorfer and Willcox, 2016; Lin et al., 2017; Kramer et al., 2017; Bistrian and Navon, 2017). In fact, several recent studies have investigated the suitability of ANN variants for fully non-intrusive dynamics capture in latent space (Mohan and Gaitonde, 2018; Wan et al., 2018; Wang et al., 2018).

With regard to the ROM methodology, we implement our framework in a context similar to the works of Narayanan et al. (1999) and Khibnik et al. (2000) where the POD-GP approach was replaced by an ANN-based evolution of the temporal dynamics of the ROM for the purpose of active separation control in a planar diffuser. Their dataset for ANN training was obtained both from direct numerical simulations of the Navier-Stokes equations as well as experiments. In Sahan et al. (1997), an ANN was used to devise a feedback control circuit for a transitional flow in a grooved channel. The dataset for ANN training was obtained through a POD-GP simulation itself. Both works demonstrated the unique advantages of the ANN approach, for example, the reduction of a system of governing PDEs to a no-equation dynamic system through data-driven learning. We note that ANNs also find great utilization in the field of feedback flow control where they are used to generate a direct mapping of flow measurements to actuator control systems (Gillies, 1995, 1998, 2001; Faller and Schreck, 1997; Hocevar et al., 2004; Efe et al., 2004, 2005; Lee et al., 1997). General machine learning based methods (of which ANN is a subset) are also growing in popularity for fluid flow based applications and represent a computationally viable alternative to the full Navier-Stokes equations (Bright et al., 2013; Gautier et al., 2015; Müller et al., 1999). An excellent source of information about some viable machine learning classes can be found in Ling and Templeton (2015). In turbulence modeling, ANN approaches have also been used in combination with reduced degree-of-freedom implementations of the Navier-Stokes framework, for instance in the quantification of errors in functional closure models (Singh and Duraisamy, 2016) and for data-driven closure modeling (Milano and Koumoutsakos, 2002; Tracey et al., 2013; Ling et al., 2016a,b; Maulik and San, 2017; Gamahara and Hattori, 2017).

This paper develops a framework for generating an ANN in order to create a very fast integrator for certain types of partial differential equation (PDE) systems (e.g., one-dimensional Burgers equation in our demonstrations). The ANN is created to represent a map for the trajectory of the POD basis coefficients evolving from time $n$ to time $n + 1$. The projection of snapshot data is used to train the ANN, i.e., make it choose the weights to best fit the data in the POD spanned reduced order space. Therefore, this model reduction approach is denoted as PON-ANN in our study. Using only snapshots of the state variables, our data-driven approach can thus be considered truly non-intrusive, since any prior information about the underlying governing equations is not required for generating the reduced order model. This idea can be considered as a competitor to the POD-GP, and the advantages are that this is an equation-free method, and also it can be more accurate since it can
be trained with more data. Although many new model order reduction techniques perform better than standard POD-GP model, as highlighted above, we focus here on comparing the proposed approach with the POD-GP model only. Future comparative studies should be carried out to explore the feasibility of such non-intrusive models. Before proceeding with our discussion, we describe the major questions we wish to answer in this work:

- Can the POD-ANN non-intrusive ROM approach be used for transient dynamical systems such as nonlinear Burgers equation?
- How does an ANN-based evolution of the modal coefficients compare to the POD-GP approach?
- Can a trained POD-ANN (using a particular dataset) be used to interpolate or extrapolate to flow problems with slight differences in temporal evolution and physics (i.e., for different control parameter values such as Reynolds number)?
- Which POD-ANN architecture is more stable and robust in forecasting beyond the training data range?

To answer these questions, we investigate a test case given by the nonlinear viscous Burgers equation to compare both the POD-ANN and POD-GP approaches.

2. Mathematical modeling

The viscous Burgers equation is used in our study as a test case to explain the proposed model reduction framework. In its original form, the Burgers equation can be written as

\[
\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \frac{1}{Re} \frac{\partial^2 u}{\partial x^2}, \quad x \in [0, 1], \quad t \in [0, 1],
\]

(1)

where \( Re \) is a non-dimensional Reynolds number. We must mention here that this equation is generally considered a framework for preliminary evaluation of numerical methods for analysis of fluid flow applications as it possesses the hallmarks of general nonlinear multi-dimensional advection-diffusion problems. We consider a convective system in our model order reduction assessments since it characterizes localized flow structures such as shock waves. The above PDE can be solved exactly to obtain an analytical formulation for time evolution of the the field variable \( u(x, t) \) given by (Maleewong and Sirisup, 2011)

\[
u(x, t) = \frac{x}{t_0 + 1} + \sqrt{\frac{t_0 + 1}{t_0}} \exp(Re \frac{x^2}{4t_0 + 1}),
\]

(2)

where \( t_0 = \exp(\frac{Re}{8}) \). This exact expression is used to generate snapshot data for our forthcoming model order reduction analysis. The database is constituted with Eq. (2) using \( N_x = 1024 \) spatial collocation points at each snapshot. Our database consists of snapshots obtained by using 10 equally spaced values of Reynolds number, i.e., \( Re = [100, 200, \ldots, 1000] \). Figure 1 shows the space-time behavior of the four representative solutions at different Reynolds number.
3. Proper orthogonal decomposition

In this section we explore the POD approach for reduced order modeling to capture unsteady, convective and non-periodic dynamics of the underlying governing partial differential equations. We elaborate the procedure for calculating the orthonormal bases and corresponding coefficients for the Burgers problem and note that an extension to higher dimensions is straightforward. For a detailed discussion on this technique we refer the reader to San and Iliescu (2014).

A representative POD basis can be constructed from any arbitrary scalar field variable \( f \) at different times (also known as snapshots). These snapshots are generally obtained by solving the governing equations we are attempting to model using a regular full-order model (FOM) approach. The numerical methods used to obtain this field data are explained in Section 6. In the following, we will utilize the index \( n \) to indicate a particular snapshot in time. For the POD approach we utilize a total of \( N \) snapshots for the field variable, i.e., \( f(x, t^{(n)}) \) for \( n = 1, 2, ..., N \). The flow field data thus obtained is first decomposed into the mean and fluctuating part as follows

\[
  f(x, t) = \bar{f}(x) + f'(x, t), \quad \bar{f}(x) = \frac{1}{N} \sum_{n=1}^{N} f(x, t^{(n)}),
\]

where \( \bar{f} \) implies a temporal averaging of a particular point value in the field and \( f' \) contains
the fluctuating quantity. We must remark here that $\bar{f}$ is a function of space alone whereas $f'$ is a function of both space and time. A correlation matrix may be constructed using the fluctuating components of the snapshots to give

$$C_{ij} = \int_\Omega f'(x, t^{(i)}) f'(x, t^{(j)}) dx,$$

where $\Omega$ is the entire spatial domain and $i$ and $j$ refer to the $i^{th}$ and $j^{th}$ snapshots. The time correlation data matrix $C$ is a non-negative symmetric square matrix of size $N \times N$. If we define the inner product of any two fields $f_1$ and $f_2$ as

$$(f_1, f_2) = \int_\Omega f_1(x) f_2(x) dx$$

we may express the correlation matrix as $C_{ij} = (f'(x, t^{(i)}), f'(x, t^{(j)}))$. In this study, we use the well-known Simpson’s 1/3 integration rule for a numerical computation of the inner products. The optimal POD basis functions are obtained by performing an eigendecomposition of the $C$ matrix. This has been shown in detail in the POD literature (see, e.g., Sirovich (1987); Holmes et al. (1998); Ravindran (2000)). The eigenvalue problem can be written in the following form:

$$CW = W \Lambda,$$

where $\Lambda = \text{diag}[\lambda_1, \lambda_2, ..., \lambda_N]$ is a diagonal matrix containing the eigenvalues of this decomposition and $W = [w^1, w^2, ..., w^N]$. The eigenvalues are stored in descending order, $\lambda_1 \geq \lambda_2 \geq ... \geq \lambda_N$. Then the orthogonal POD basis functions can be written as

$$\phi_1(x) = \sum_{n=1}^{N} w^1_n f'(x, t^{(n)}), \quad \phi_2(x) = \sum_{n=1}^{N} w^2_n f'(x, t^{(n)}), \quad ..., \quad \phi_N(x) = \sum_{n=1}^{N} w^N_n f'(x, t^{(n)}),$$

where $w^k_n$ is the $n^{th}$ component of the $k^{th}$ eigenvector $w^k$. Therefore, we emphasize that the POD modes are ranked according to the magnitude of their eigenvalue. The eigenvectors must also be normalized in order to satisfy the condition of orthonormality between bases:

$$(\phi_i, \phi_j) = \begin{cases} 1, & i = j \\ 0, & i \neq j. \end{cases}$$

It can be shown that, for Eq. (8) to be true, the eigenvector $w^k$ must satisfy the following equation:

$$\sum_{n=1}^{N} w^k_n w^k_n = \frac{1}{\lambda_k},$$

where $\lambda_k$ is the $k^{th}$ eigenvalue associated to the eigenvector $w^k$. In practice, most of the subroutines for solving the eigensystem given in Eq. (6) return the eigenvector matrix $W$. The time correlation data matrix $C$ is a non-negative symmetric square matrix of size $N \times N$. If we define the inner product of any two fields $f_1$ and $f_2$ as

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$$CW = W \Lambda,$$
having all the eigenvectors normalized to unity. In that case, the orthogonal POD bases are given by

\[ \phi_k(x) = \frac{1}{\sqrt{\lambda_k}} \sum_{n=1}^{N} w_k^n f'(x, t^n) \]  

(10)

where \( \phi_k(x) \) is the \( k \)th POD basis function. We remark here that the POD should be carried out independently for each field variable for multidimensional governing equations. We collect 101 snapshots between time \( t = 0 \) and \( t = 1 \) at each Reynolds number between \( Re = 100 \) and \( Re = 1000 \) using 10 different realizations (i.e., \( Re = [100, 200, ..., 1000] \)). Figure 2 shows the eigenvalues of the time correlation matrix \( C \) and the corresponding relative information content (RIC) index, which is defined as (Gunzburger, 2012; Tallet et al., 2016)

\[ RIC_k = \frac{\sum_{n=1}^{k} \lambda_n}{\sum_{n=1}^{N} \lambda_n} \times 100, \]  

(11)

where the number of snapshots is set to \( N = 1010 \) in our study. As shown in Table 1, the RIC index measures the percentage of captured energy for different number of POD modes. It is evident that a small number of modes are enough to capture the spatial and temporal dynamics of the system well. Therefore, we present our ROM analyses using \( R = 6 \) largest (most energetic) POD modes capturing approximately 98.27 % of the energy. Figure 3 illustrates the associated POD basis functions utilized in this study.

Figure 2: Eigenvalues of the snapshot correlation matrix (left) and the corresponding relative information content (RIC) index (right) for the space-time solution of the viscous Burgers equation. Note that we retain only the most energetic 6 POD modes in our model order reduction analysis.
Table 1: The relative information content (RIC) index that measures the percentage of captured energy for different number of POD modes.

| R   | RIC(%)   |
|-----|----------|
| 1   | 70.0273  |
| 2   | 87.1354  |
| 3   | 93.0921  |
| 4   | 95.9237  |
| 5   | 97.3327  |
| 6   | 98.2723  |
| 10  | 99.6722  |
| 20  | 99.9897  |

4. Galerkin projection methodology

The POD basis functions account for the essential dynamics of the underlying governing equations. After these data-driven empirical POD basis functions are obtained, a set of nonlinear ODEs can be derived by using a Galerkin projection. To build a projection based ROM, we truncate the system by considering the first $R$ largest POD basis functions with $R \ll N$. These POD modes correspond to the $R$ largest eigenvalues, $\lambda_1, \lambda_2, \ldots, \lambda_R$. Using these first $R$ largest POD basis functions the field variables are then approximated.

Figure 2 shows a collection of the modal values (i.e., the eigenvalues) of the POD-ROM analysis. The main motivation behind the construction of a POD-GP is evident when it is observed that the 6 most energetic modes retain 98.27% of the total energy of the system. We note that these bases are all orthogonal to each other. For our test case, i.e., $u(x, t) = \bar{u} + u'(x, t)$, the Galerkin projection can be carried out in the following manner

$$u'(x, t) = \sum_{k=1}^{R} a_k(t)\phi_k(x),$$  \hspace{1cm} (12)

where $a_k$ are the time dependent coefficients, and $\phi_k$ are the space dependent modes. To derive the POD-ROM, we first rewrite the Burgers equation (i.e., Eq. (1)) in the following form

$$\frac{\partial u}{\partial t} = L[u] + N[u; u],$$  \hspace{1cm} (13)

where $L[f] = \frac{1}{Re} \frac{\partial^2 f}{\partial x^2}$ is the linear operator and $N[f; g] = -f \frac{\partial g}{\partial x}$ is the nonlinear operator. We emphasize that the procedure can be easily extended to the general nonlinear PDEs (e.g., Boussinesq or Navier-Stokes equations). By applying this projection to our nonlinear system (i.e., multiplying Eq. (13) with the basis functions and integrating over the domain), we obtain the Galerkin POD-ROM, denoted POD-GP:

$$\left( \frac{\partial u}{\partial t}, \phi_k \right) = (L[u], \phi_k) + (N[u; u], \phi_k), \text{ for } k = 1, 2, \ldots, R.$$  \hspace{1cm} (14)
Substituting Eq. (12) into Eq. (14), and simplifying the resulting equation by using the condition of orthonormality given in Eq. (8), the POD-GP implementation can be written as follows:

$$\frac{da_k}{dt} = B_k + \sum_{i=1}^{R} S_{ik} a_i + \sum_{i=1}^{R} \sum_{j=1}^{R} N_{ijk} a_i a_j, \quad \text{for} \quad k = 1, 2, \ldots, R,$$

(15)
where

\[
B_k = (L[\bar{u}] + N[\bar{u}; \bar{u}], \phi_k), \quad (16)
\]

\[
\Sigma_{ik} = (L[\phi_i] + N[\bar{u}; \phi_i] + N[\phi_i; \bar{u}], \phi_k), \quad (17)
\]

\[
N_{ijk} = (N[\phi_i; \phi_j], \phi_k). \quad (18)
\]

The POD-GP given by Eq. (15) consists of \( R \) coupled ODEs and can be solved by a standard numerical method (such as the third-order Runge-Kutta scheme that was used in this study). The number of degrees of freedom of the system is now significantly lower. The vectors, matrices and tensors in Eqs. (16)-(18) are also precomputed quantities, which results in a dynamical system that can be solved very efficiently. To complete the dynamical system given by Eq. (15), the initial condition is given by using the following projection:

\[
a_k(t = 0) = (u(x, t = 0) - \bar{u}(x), \phi_k), \quad (19)
\]

where \( u(x, t = 0) \) is the physical initial condition of the problem given in Eq. (2).

5. Artificial neural networks

A simple feed-forward artificial neural network consists of \( L \) layers with each layer possessing a predefined number of unit cells called neurons and is utilized for establishing nonlinear maps between two spaces of potentially different dimensionality through a supervised learning. Each of these layers may be considered an intermediate step in successive transformations between the input and output space and has an associated transfer function and each unit cell has an associated bias. Therefore, in our choice of network, any input to the neuron has a bias added to it followed by an activation through the transfer function. To describe this process using equations, we have

\[
s_l^i = \sum_j W_{lj}^i X_{l-1}^j, \quad (20)
\]

as the state of the signal in the \( l^{th} \) layer receiving a set of inputs from the \((l-1)^{th}\) layer, where \( W_{lj}^i \) stands for a matrix of weights linking the \( l-1 \) and \( l \) layers with \( X_{l-1}^j \) being the output of the \((l-1)^{th}\) layer. The output signal of the \( l^{th} \) layer is now given by

\[
X_l^i = F(s_l^i + b_l^i), \quad (21)
\]

where \( b_l^i \) is the biasing parameter for the \( i^{th} \) neuron in the \( l^{th} \) layer. Every node (or unit cell) has an associated transfer function which acts on its input and bias to produce an output which is ‘fed forward’ in the network. The nodes which take the raw inputs of our training data set (i.e., the nodes of the first layer in the network) perform no computation (i.e., they do not have any biasing or activation through a transfer function). The next layers are a series of unit cells which have an associated bias and activation function which perform computation on their inputs. These are called the hidden layers due to the indeterminate
nature of their mathematical operations. The final layer in the network is that of the outputs. The output layer generally has a linear activation function with a bias which implies a simple summation of inputs incident to a unit cell with its associated bias. A mathematical description of this mapping operation can be expressed as

\[ M: \{p_1, p_2, \ldots, p_P\} \in \mathbb{R}^P \rightarrow \{q_1, q_2, \ldots, q_Q\} \in \mathbb{R}^Q. \] (22)

where \( P \) and \( Q \) are the dimensions of the input and output spaces. In this investigation, we have used one hidden layer of neurons between the set of inputs and targets with a tan-sigmoid activation function. The tan-sigmoid function can be expressed as

\[ F(a) = \frac{2}{1 + \exp(-2a)} - 1. \] (23)

The transfer function \( F \) calculates the neuron’s output given its net input. In theory, any differentiable function can qualify as an activation function (Zhang et al., 1998), however, only a small number of functions which are bounded, monotonically increasing and differentiable are used for this purpose. The choice of using ANN is motivated by its excellent performance as a forecasting tool (Dawson and Wilby, 1998; Kim and Valdés, 2003) and its general suitability in the machine learning and function estimation domain (e.g., see Haykin et al. (2009) and references therein).

5.1. POD transforms

The key idea in our learning problem is the model reduction. Instead of posing our learning problem in a full-order space, we first apply projection to the full-order data and generate coefficients on a reduced subspace, and then construct the learning on this reduced-order space. Here we briefly describe the main procedure to generate such non-intrusive reduced order models. Let’s consider we have snapshots either coming from the full-order numerical simulations or experiments. Section 3 illustrates how to obtain representative POD modes. At any snapshot, \( u^{(n)} \), our encoder operator can be written using the definition of inner product as

\[ a_k^{(n)} = \langle u^{(n)} - \bar{u}, \phi_k \rangle, \quad \text{for} \quad k = 1, 2, \ldots, R, \] (24)

where basis functions \( \phi_k \) and the mean field \( \bar{u} \) are already available. This can be interpreted as a forward POD transform from the FOM space to the ROM space. We then perform our supervised learning for \( a_k^{(n)} \) to train a map for computing the modal coefficient trajectories. The details of the neural network architectures can be found in more detail in the following section. Finally, the decoder operator is defined as

\[ u^{(n)} = \bar{u} + \sum_{k=1}^{R} a_k^{(n)} \phi_k, \] (25)

which can be referred to as an inverse POD transform. Indeed this expression can be used to generate full-order space data from the modal coefficients at any time.
5.2. ANN training

The training of a desired ANN is carried out by minimizing the error between the target and the inputs to determine a set of best fit parameters. This is also known as supervised learning where labeled data is utilized for gradient based optimization. Our best fit parameters obtained from optimization are biases and linear weights that have captured the underlying relationship between the targets and inputs and may now be used to predict target data for inputs a-posteriori. The main advantage of using the ANN approach over traditional statistical regression models is that comparatively smaller data sets for training are suitable. We shall describe the individual features of the ANN architecture relevant to our test case in the sections that follow.

The training of an ANN architecture may be undertaken by a variety of optimization algorithms. For our purpose, we have employed the Bayesian regularization minimization algorithm which is one of the popular approaches to ANN training within the framework of the MATLAB `nnstart` toolbox and particularly useful for noisy data ( Gençay and Qi, 2001). We remark that other training methods may also be used such as Levenberg-Marquardt (Levenberg, 1944) and nature-inspired heuristic algorithms such as particle swarm (Eberhart and Kennedy, 1995), differential evolution (Storn and Price, 1997) or leapfrogging (Rhinehart et al., 2012). However, access to gradient information (through the use of a differentiable cost-function) heavily favors the use of gradient-based optimizers. The process of backpropagation, where the errors in the output of the network are propagated back through the network to determine optimal search directions in parameter space is well suited for our particular architecture which utilizes solely one hidden layer. However, a detailed hyperparamater search is beyond the scope of this study since we would like to focus on the performance assessment of the proposed architectures using default options available in `nnstart` (e.g., mean-squared error loss function).

For the purpose of tracking the training of the ANN and evaluating its performance, we utilize three different subsets of our overall data with 70% being utilized for training, and 15% being utilized for validation and testing each. The training data set is used to adjust the weights of the ANN. The validation data set is used to minimize overfitting, i.e., it is used to verify whether an increase in accuracy over the training data set also yields an increased accuracy over a data set that is not revealed to the network a-priori. The validation data set is used for testing the final solution of the neural network. We note that the input and output features are all scaled between their minimum and maximum values and the scaling parameters are ported for deployment. Within the POD context, the trained ANN can be used as a ‘blackbox’ which takes the present state of the latent space dynamics as the input along with a dimensionless parameter and the present time and uses the memory of its training to determine the time evolution. Although not shown here for brevity, we note that the regression analysis shows a perfect correlation in the time history of the temporal modes and develops a system which could predict their evolution accurately.

5.3. ANN architecture for model order reduction

In the present study, we devise a framework for developing an ANN in order to create a very fast (equation-free or non-intrusive) integrator/solver for certain types of PDE systems.
This idea can be considered as a competitor for POD-GP, with the advantage this is an equation-free and purely data-driven with associated benefits in computational efficiency and model-form uncertainty reduction.

Figure 4 illustrates the ANN architectures utilized in this study. In these architectures, several inputs are provided, out of which the majority are basis coefficients at the current timestep in the POD-ROM evolution and the remaining two are Reynolds number and time. The training data is also arranged in the same manner. Our outputs are given by just the coefficient values required for a field reconstruction in the next step. The figure referred to above shows the number of inputs (modal POD coefficients, Reynolds number and time) and the number of outputs (modal POD coefficients at the next timestep). The middle layer l contains the neurons with the tan-sigmoid activation function and these are varied in our investigation. Therefore our first mapping approach, sequential network (SN) architecture, equipped by $R = 6$ modes can be given by

$$ M_1 : \{ \text{Re}, t^{(n)}, a_1^{(n)}, a_2^{(n)}, \ldots, a_6^{(n)} \} \in \mathbb{R}^8 \rightarrow \{a_1^{(n+1)}, a_2^{(n+1)}, \ldots, a_6^{(n+1)}\} \in \mathbb{R}^6. \quad (26) $$

We denote this model as POD-ANN-SN. As discussed by Chen et al. (2018a), machine learning models such as residual networks (He et al., 2016) define a discrete sequence of finite transformations to a hidden state in the following form (Haber and Ruthotto, 2017; Lu et al., 2017):

$$ a_k^{(n+1)} = a_k^{(n)} + r_k, \quad \text{for} \quad k = 1, 2, \ldots, R, \quad (27) $$

where $r_k$ refers to the residual that we can approximate using an ANN. This can be seen as an Euler update in a discrete sense as follows (i.e., defining the residual as $r_k = \Delta t f_k$)

$$ a_k^{(n+1)} = a_k^{(n)} + \Delta t f_k, \quad (28) $$

where $\Delta t = t^{(n+1)} - t^{(n)}$, and $f_k$ is the slope

$$ f_k = \frac{a_k^{(n+1)} - a_k^{(n)}}{\Delta t}. \quad (29) $$

To investigate the performance of such a residual network, a similar network shown also in Figure 4 is also designed for $R = 6$ largest POD modes, which can be defined as

$$ M_2 : \{ \text{Re}, t^{(n)}, a_1^{(n)}, a_2^{(n)}, \ldots, a_6^{(n)} \} \in \mathbb{R}^8 \rightarrow \{r_1, r_2, \ldots, r_6\} \in \mathbb{R}^6. \quad (30) $$

and the POD coefficients are then obtained by Eq. (27) before the consequent step. This model is denoted as POD-ANN-RN in our study. Details related to data acquisition for our network training are provided in Section 7. We highlight that the main objective of the present work is to test and evaluate the characteristics of the proposed POD-ANN approaches and compare them with the true solution obtained by a projection to the full order model (FOM). Our assessments also includes intrusive results generated by the standard POD-GP framework.
6. Numerical methods

In this section, we provide a brief description of the numerical methods employed in this study. First of all, we implement compact difference schemes to approximate the differential operators (both linear and nonlinear terms) in the POD-GP. Compact difference schemes have shown their ability to reach the objectives of high-accuracy and low computational cost for many problems in fluid dynamics (Lele 1992; Wang and Liu 2002).

In a compact difference scheme, the first order derivatives can be computed accordingly (Lele 1992)

\[ \alpha f_i' - f_{i-1}' + \alpha f_{i+1}' = a \frac{f_{i+1} - f_{i-1}}{2\Delta x} + b \frac{f_{i+2} - 2f_i + f_{i-2}}{4\Delta x}, \]  

which gives rise to an \( \alpha \)-family of tridiagonal schemes with \( a = \frac{2}{3}(\alpha + 2) \), and \( b = \frac{1}{3}(4\alpha - 1) \). The subscript \( i \) represents the spatial grid index in the considered direction, and \( \Delta x \) is the uniform grid spacing in that direction. Here, \( \alpha = 0 \) leads to the explicit non-compact fourth-order scheme for first derivative computations. A classical compact fourth-order scheme, which is also known as Padé scheme, is obtained by setting \( \alpha = 1/4 \), resulting in \( b = 0 \) with a compact dependence. In computing the linear diffusion term, the second derivative compact scheme is given by

\[ \alpha f''_{i-1} + f'_i + \alpha f''_{i+1} = a \frac{f_{i+1} - 2f_i + f_{i-1}}{\Delta x^2} + b \frac{f_{i+2} - 2f_i + f_{i-2}}{4\Delta x^2}, \]  

Figure 4: The non-intrusive model order reduction enablers for transient flows with the sequential network \( \mathcal{M}_1 \) (left) and residual network \( \mathcal{M}_2 \) (right) ANN architectures. We illustrate our architectures for one physical control parameter \( Re \) and use only one hidden layer equipped with \( M \) neurons and employing \( R \) most energetic POD modes.
where \( a = \frac{4}{3}(1 - \alpha) \), and \( b = \frac{1}{3}(10\alpha - 1) \). For \( \alpha = 1/10 \) the classical fourth-order Padé scheme is obtained. The high-order one-sided derivative formulas are used for the Dirichlet boundary conditions to complete a tridiagonal system of equations for both the first and second order derivative formulas (Carpenter et al., 1993).

As a time integrator, we use an optimal third-order accurate total variation diminishing Runge-Kutta (TVDRK3) scheme that is given as (Gottlieb and Shu, 1998)

\[
\begin{align*}
    u_i^{(1)} &= u_i^{(n)} + \Delta t \mathcal{L}(u_i^{(n)}), \\
    u_i^{(2)} &= \frac{3}{4} u_i^{(n)} + \frac{1}{4} u_i^{(1)} + \frac{1}{4} \Delta t \mathcal{L}(u_i^{(1)}), \\
    u_i^{(n+1)} &= \frac{1}{3} u_i^{(n)} + \frac{2}{3} u_i^{(2)} + \frac{2}{3} \Delta t \mathcal{L}(u_i^{(2)}),
\end{align*}
\]

(33)

where \( \mathcal{L}(u) \) is the discrete operator of spatial derivatives including linear and nonlinear terms. To perform the inner products defined by Eq. (5), we compute the integral of any arbitrary function \( g(x) \) over the domain \( \Omega \) by using the Simpson’s 1/3 rule (Hoffman and Frankel, 2001)

\[
\int_\Omega g(x) dx = \frac{\Delta x}{3} \sum_{i=1}^{N_x/2-1} \left( g_{2i} + 4g_{2i+1} + g_{2i+2} \right),
\]

(34)

where \( N_x \) is the total number of grid point (e.g., an even number in order to be consistent with the Simpson’s integration rule).

7. Results

In the following we provide modal coefficient evolution data for our test case and make assessments of the relative performance of the POD-ANN and POD-GP approaches. Although measuring the state evolution error is not equivalent to measuring the error in any standard function space norm or quantity of interest, we found that it simplifies our demonstration and presents an efficient way in our assessments between the proposed approaches and standard POD-GP model. In addition, we also provide details related to the calculation of POD bases and their coefficients for the use of data-driven learning and a-posteriori deployment.

POD bases for the viscous Burgers equation problem are generated by accumulation of snapshot data from the exact solution given by Eq. (2). As described in Section 3, snapshot sampling is done at different Reynolds numbers between 100 and 1000 at intervals of 100 each (i.e., \( Re = 100, Re = 200, ..., Re = 1000 \)). Each choice of a Reynolds number is sampled for 101 snapshots in time between \( t = 0 \) and \( t = 1 \). This leads to a total of 1010 snapshots to construct a snapshot correlation data matrix \( C \) with the size of \( 1010 \times 1010 \). Although we generate such 1010 snapshots from the closed form expression, we remark that these snapshots can be obtained from experiments or full order model numerical simulations in most real-life problems. The physics of these multiple snapshots for each control parameter is thus spanned by a single set of optimal bases used for subsequent ANN training or Galerkin projection. The ANN training data has been reconstructed with a timestep of \( \Delta t = 0.001 \) for
the same database with 10 different $Re$ from $t = 0$ to $t = 1$, which leads to 10,010 samples in our training set. We note that the timestep for POD-GP and POD-ANN deployment are kept identical in our numerical assessments. We highlight that the training of ANN to bypass the Galerkin procedure is an offline cost which utilizes approximately 45 seconds.

We note that the results shown in this subsection are for Reynolds number values which were not within the training dataset range for the ANN. Throughout our study we use a single-layer neural network with $M = 10$ neurons and employing $R = 6$ POD modes. Figure 5 shows a comparison at $Re = 750$ on predicting the POD basis coefficients. Our numerical assessment with Reynolds number of 750 represents the interpolative abilities of both POD-GP and POD-ANN approaches. It is clear that the POD-ANN approach preserves state accuracy for a greater duration as compared to the POD-GP approach. At this $Re$, the performance of the both POD-ANN approaches is remarkable. Indeed for the first mode, the POD-ANN approaches preserve the amplitude and phase behavior for the coefficient very well. Coefficients for higher modes tend to show a slight variation in amplitude when compared to DNS data but this variation is quite negligible when compared along the POD-GP approach which shows extensive loss of phase and amplitude at higher modes.

Our next test at $Re = 1250$ represents the extrapolative abilities of both POD-GP and POD-ANN models, in which the training data is generated from snapshots between $Re = 100$ and $Re = 1000$. Figure 6 shows a similar study where we have retained only 6 POD modes to predict the flow field at higher $Re$. It can be easily seen that both non-intrusive POD-ANN-SN and POD-ANN-RN approaches yield stable and accurate results. Furthermore, one can immediately notice a significant inaccuracy in the POD-GP approach. Considering only six largest POD modes might not be rich enough to represent the underlying transient behavior. However, the POD-ANN approaches still perform well in terms of capturing the amplitude of the first six modes even though slight variations in the phase for the highest modes are seen in the $Re = 1250$ case as shown in Figure 6. Although the POD-GP method shows a considerable deviation from the truth as compared to the POD-ANN approaches, when examining the difference between the POD-ANN-SN and POD-ANN-RN non-intrusive methodologies, we can see that the POD-ANN-RN provides more accurate estimates as compared to the POD-ANN-SN approach which does a decent job at capturing amplitude and phase behavior at this Reynolds number.

To further investigate the behavior of the proposed ANN architectures at higher Reynolds number, we test the proposed non-intrusive frameworks for $Re = 2000$. To this end, Figure 7 presents our numerical assessments at $Re = 2000$ that is far beyond the parameter range utilized in the training database. This examination points out that the POD-ANN-RN is more stable than the POD-ANN-SN. This figure also demonstrates significantly more accurate behavior shown by the POD-ANN-RN approach. In terms of computational costs, the POD-ANN approaches equipped by $M = 10$ neurons do not require the use of Runge-Kutta based time stepping and is thus observed to be two orders of magnitude faster than the POD-GP. However, we note that the running time is significantly less than a second for both non-intrusive POD-ANN and intrusive POD-GP approaches with $R = 6$ modes. Indeed, we observe that the computational performance of both POD-GP and POD-ANN
Figure 5: Numerical assessments of the POD-GP and POD-ANN models with $R = 6$ POD modes applied to the Burgers problem for an interpolatory out-of-sample parameter $Re = 750$, where the training set includes snapshots between $Re = 100$ and $Re = 1000$. The non-intrusive POD-ANN-SN and POD-ANN-RN results are obtained using $M = 10$ neurons in the hidden layer. Note that the exact solution projected to the reduced order space is shown with the black solid line with circles.

...codes is dominated by data-writing for post-processing.

We stress here that bypassing the computation of the reduced system obtained through Galerkin projection and truncation ensures that our non-intrusive frameworks are purely data-driven. The non-intrusive nature of the POD-ANN methodology and the lack of a system of equations that needs to be solved for discretely in time thus leads to an attractive alternative to the POD-GP implementation. Although we illustrate that a robust surrogate model can be obtained by a residual network based ANN architecture without a need to
Figure 6: Numerical assessments of the POD-GP and POD-ANN models with $R = 6$ POD modes applied to the Burgers problem at $Re = 1250$, where the training set includes snapshots between $Re = 100$ and $Re = 1000$. The non-intrusive POD-ANN-SN and POD-ANN-RN results are obtained using $M = 10$ neurons in the hidden layer. Note that the exact solution projected to the reduced order space is shown with the black solid line with circles.

access to the governing equations, we highlight that the intrusive or any other models generated from the first principles allow one to prove important properties like numerical stability, convergence, energy conservation, etc., which are all seemingly lost in blackbox ANN models. However, we refer the readers to Xu and Darve (2019) for an excellent discussion of the key issues in utilizing neural network approaches to inverse problems in differential equations such as universal approximation, regularization, curse of dimensionality and computational efficiency.
Figure 7: Numerical assessments of the POD-GP and POD-ANN models with $R = 6$ POD modes applied to the Burgers problem at $Re = 2000$, where the training set includes snapshots between $Re = 100$ and $Re = 1000$. The non-intrusive POD-ANN-SN and POD-ANN-RN results are obtained using $M = 10$ neurons in the hidden layer. Note that the exact solution projected to the reduced order space is shown with the black solid line with circles.

8. Concluding remarks

In this study, we have investigated the predictive performance of two different POD-enabled model order reduction approaches for the time evolution of the state. First, we look at the classical Galerkin projection approach where a system of ODEs is developed for the evolution of each state coefficient (corresponding to a scalar field variable). Second, we have developed an equation-free artificial neural network (ANN) framework through data-driven learning which predicts the evolution of the modes without the explicit time stepping
approach used in the Galerkin projection. We give a demonstration of the proposed ANN framework for two architectures: (i) a sequential network (SN) where we compute the trajectory of POD coefficients directly, and (ii) a residual network (RN) where we train our model with a discrete residual information, rather than using the sequential state variables. Our numerical assessments are presented for the viscous Burgers equation benchmark problem at various Reynolds number. We notice that the POD-ANN approaches outperforms POD-GP consistently across the different testing environments and proves excellent in both interpolating and extrapolating its learning from obtained data sets for simulations with different control parameters and underlying physics. Furthermore, we demonstrate that POD-ANN-RN yields significantly more stable results especially for the extrapolatory parameter range that are outside of the training range. We revisit the questions we posed in the introduction to conclude our investigation:

- We can conclude with confidence that the POD-ANN approach is superior to the POD-GP approach for the Burgers equation especially for highly truncated systems (i.e., systems represented by a few modes only).

- We can also state that POD-ANN is a viable tool for extrapolation and interpolation beyond the data sets used to train its learning. While it is expected to provide good results within the bounds of the training data set, the Burgers case detailed its capacity to predict the coefficient evolution of different Reynolds numbers in an accurate manner with much greater accuracy than POD-GP.

- We find that the residual network formulation is significantly more stable than the sequential network formulation.

Furthermore, we present the following observations which were made during this study:

- Once trained, the POD-ANN can be considered to be an ‘equation-free’ ROM architecture.

- Although the POD-GP gives us usually better approximations while increasing the retained number modes, the POD-ANN is better suited to more radically truncated systems since an ANN with higher degrees of freedom may show statistical errors on time evolution.

- POD-ANN represents a viable framework for modeling highly nonlinear systems when data is available incrementally and constant updates are required. It is also very well suited to situations where model interpretability is not a major concern.

- While our training algorithm utilized the standard Bayesian regularization algorithm, it must be noted that there exists many training methodologies which may better preserve the underlying statistical relationships of the inputs and outputs. The effect of training approaches is a future direction of study for this work.
We also recognize that the POD-ANN, despite its obvious benefits, also leads to the drawback of reduced interpretability due to its blackbox nature. In addition, the a-priori requirement of training implies poor generalizability of the POD-ANN framework to completely different physical regimes.

Additionally we note that the use of more advanced regularization techniques in the ANN training (an aspect not considered in this study) would potentially lead to more stable non-intrusive ROM methodologies.

Training with noisy data may also be studied for obviating the need for closure modeling in the POD-GP framework.

Table 2: A comparison of intrusive and non-intrusive methodologies describing strengths and weaknesses for each framework.

| Intrusive ROMs (POD-GP)                        | Non-intrusive ROMs (POD-ANN)                  |
|-----------------------------------------------|----------------------------------------------|
| + Solid foundations based on physics and first principles (high interpretability) | - Mostly blackbox solvers and packages (low interpretability) |
| - Need access to the governing equations       | + Only need observed/measured data           |
| + Generalizes well to new problems with similar physics | - Poor generalization on unseen problems      |
| - Implementations are mostly problem dependent | + Extremely convenient for tensor based computing |
| + Numerical analysis is easier                 | - Hard to establish rules for its convergence, stability, and energy conservation |

An overview of our conclusions are provided in Table 2 which contrasts the difference between POD-GP and POD-ANN type methods for ROMs. Although our study utilizes a shallow neural network architecture, we believe that the integration of machine learning tools such as deep learning into the mainstream reduced order modeling ideology would represent a significant increase in its viability for both theoretical and application oriented goals. We would like to conclude that POD-ANN-RN would be a viable ROM procedure for convective flows where the response of complex system dynamics is desired to be represented accurately by a few representative modes.

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References

Ahmed, M., San, O., 2018. Stabilized principal interval decomposition method for model reduction of nonlinear convective systems with moving shocks. Computational and Applied Mathematics 37 (5), 6870–6902.

Akhtar, I., Borggaard, J., Burns, J. A., Imtiaz, H., Zietsman, L., 2015. Using functional gains for effective sensor location in flow control: a reduced-order modeling approach. Journal of Fluid Mechanics 781, 622–656.

Amsaleg, D., Farhat, C., 2012. Stabilization of projection-based reduced-order models. International Journal for Numerical Methods in Engineering 91 (4), 358–377.

Aubry, N., Holmes, P., Lumley, J. L., Stone, E., 1988. The dynamics of coherent structures in the wall region of a turbulent boundary layer. Journal of Fluid Mechanics 192 (1), 115–173.

Audouze, C., De Vuyst, F., Nair, P. B., 2013. Nonintrusive reduced-order modeling of parametrized time-dependent partial differential equations. Numerical Methods for Partial Differential Equations 29 (5), 1587–1628.

Benner, P., Gugercin, S., Willcox, K., 2015. A survey of projection-based model reduction methods for parametric dynamical systems. SIAM Review 57 (4), 483–531.

Berkooz, G., Holmes, P., Lumley, J. L., 1993. The proper orthogonal decomposition in the analysis of turbulent flows. Annual Review of Fluid Mechanics 25 (1), 539–575.

Bistrian, D., Navon, I., 2015. An improved algorithm for the shallow water equations model reduction: dynamic mode decomposition vs POD. International Journal for Numerical Methods in Fluids 78 (9), 552–580.

Bistrian, D. A., Navon, I. M., 2017. Randomized dynamic mode decomposition for nonintrusive reduced order modelling. International Journal for Numerical Methods in Engineering 112 (1), 3–25.

Borggaard, J., Hay, A., Pelletier, D., 2007. Interval-based reduced order models for unsteady fluid flow. International Journal of Numerical Analysis & Modeling 4 (3-4), 353–367.

Borggaard, J., Iliescu, T., Wang, Z., 2011. Artificial viscosity proper orthogonal decomposition. Mathematical and Computer Modelling 53 (1-2), 269–279.

Bright, I., Lin, G., Kutz, J. N., 2013. Compressive sensing based machine learning strategy for characterizing the flow around a cylinder with limited pressure measurements. Physics of Fluids 25 (12), 127102.

Brunton, S. L., Noack, B. R., 2015. Closed-loop turbulence control: progress and challenges. Applied Mechanics Reviews 67 (5), 050801.

Buffoni, M., Camarri, S., Iollo, A., Salvetti, M. V., 2006. Low-dimensional modelling of a confined three-dimensional wake flow. Journal of Fluid Mechanics 569, 141–150.

Bui-Thanh, T., Willcox, K., Ghattas, O., van Bloemen Waanders, B., 2007. Goal-oriented, model-constrained optimization for reduction of large-scale systems. Journal of Computational Physics 224 (2), 880–896.

Carlberg, K., Farhat, C., 2011. A low-cost, goal-oriented ‘compact proper orthogonal decomposition’ basis for model reduction of static systems. International Journal for Numerical Methods in Engineering 86 (3), 381–402.

Carpenter, M. H., Gottlieb, D., Abarbanel, S., 1993. Stable and accurate boundary treatments for compact, high-order finite-difference schemes. Applied Numerical Mathematics 12 (1), 55–87.
Chen, T. Q., Rubanova, Y., Bettencourt, J., Duvenaud, D. K., 2018a. Neural ordinary differential equations. In: Advances in Neural Information Processing Systems. pp. 6571–6583.

Chen, W., Hesthaven, J. S., Junqiang, B., Qiu, Y., Yang, Z., Tihao, Y., 2018b. Greedy nonintrusive reduced order model for fluid dynamics. AIAA Journal 56 (12), 4927–4943.

Cizmas, P. G., Richardson, B. R., Brenner, T. A., O’Brien, T. J., Breault, R. W., 2008. Acceleration techniques for reduced-order models based on proper orthogonal decomposition. Journal of Computational Physics 227 (16), 7791–7812.

Cordier, L., Noack, B. R., Tissot, G., Lehnasch, G., Delville, J., Balajewicz, M., Daviller, G., Niven, R. K., 2013. Identification strategies for model-based control. Experiments in Fluids 54 (8), 1580.

Dawson, C. W., Wilby, R., 1998. An artificial neural network approach to rainfall-runoff modelling. Hydrological Sciences Journal 43 (1), 47–66.

Demuth, H. B., Beale, M. H., De Jess, O., Hagan, M. T., 2014. Neural network design. Martin Hagan, OK, USA.

Dyke, S., Spencer Jr, B., Sain, M., Carlson, J., 1996. Modeling and control of magnetorheological dampers for seismic response reduction. Smart Materials and Structures 5 (5), 565.

Eberhart, R., Kennedy, J., 1995. A new optimizer using particle swarm theory. In: Proceedings of the Sixth International Symposium on Micro Machine and Human Science. IEEE, pp. 39–43.

Efe, M., Debiasi, M., Yan, P., Ozbay, H., Samimy, M., 2005. Control of subsonic cavity flows by neural networks-analytical models and experimental validation. In: 43rd AIAA Aerospace Sciences Meeting and Exhibit. p. 294.

Efe, M. O., Debiasi, M., Ozbay, H., Samimy, M., 2004. Modeling of subsonic cavity flows by neural networks. In: Proceedings of the IEEE International Conference on Mechatronics. IEEE, pp. 560–565.

El Majd, B. A., Cordier, L., 2016. New regularization method for calibrated POD reduced-order models. Mathematical Modelling and Analysis 21 (1), 47–62.

Faller, W. E., Schreck, S. J., 1997. Unsteady fluid mechanics applications of neural networks. Journal of Aircraft 34 (1), 48–55.

Fang, F., Pain, C., Navon, I., Gorman, G., Piggott, M., Allison, P., Farrell, P., Goddard, A., 2009. A POD reduced order unstructured mesh ocean modelling method for moderate Reynolds number flows. Ocean Modelling 28 (1), 127–136.

Fortuna, L., Numari, G., Gallo, A., 2012. Model order reduction techniques with applications in electrical engineering. Springer Science & Business Media.

Freund, R. W., 1999. Reduced-order modeling techniques based on Krylov subspaces and their use in circuit simulation. In: Applied and computational control, signals, and circuits. Springer, pp. 435–498.

Gamahara, M., Hattori, Y., 2017. Searching for turbulence models by artificial neural network. Physical Review Fluids 2 (5), 054604.

Gautier, N., Aider, J.-L., Duriez, T., Noack, B., Segond, M., Abel, M., 2015. Closed-loop separation control using machine learning. Journal of Fluid Mechanics 770, 442–457.

Gençay, R., Qi, M., 2001. Pricing and hedging derivative securities with neural networks: Bayesian regularization, early stopping, and bagging. IEEE Transactions on Neural Networks 12 (4), 726–734.

Gillies, E., 1995. Low-dimensional characterization and control of non-linear wake flows. Ph.D. thesis, PhD Dissertation, University of Glasgow, Scotland.

Gillies, E., 1998. Low-dimensional control of the circular cylinder wake. Journal of Fluid Mechanics 371, 157–178.

Gillies, E., 2001. Multiple sensor control of vortex shedding. AIAA Journal 39 (4), 748–750.

Gottlieb, S., Shu, C. W., 1998. Total variation diminishing Runge-Kutta schemes. Mathematics of Computation 67 (221), 73–85.

Gunzburger, M. D., 2012. Flow control. Springer, Berlin.

Haber, E., Ruthotto, L., 2017. Stable architectures for deep neural networks. Inverse Problems 34 (1), 014004.

Hampton, J., Fairbanks, H. R., Narayan, A., Doostan, A., 2018. Practical error bounds for a non-intrusive bi-fidelity approach to parametric/stochastic model reduction. Journal of Computational Physics 368,
Haykin, S. S., Haykin, S. S., Haykin, S. S., Haykin, S. S., 2009. Neural networks and learning machines. Vol. 3. Pearson Upper Saddle River, NJ, USA.

He, K., Zhang, X., Ren, S., Sun, J., 2016. Deep residual learning for image recognition. In: Proceedings of the IEEE conference on computer vision and pattern recognition. pp. 770–778.

Hesthaven, J. S., Ubbiali, S., 2018. Non-intrusive reduced order modeling of nonlinear problems using neural networks. Journal of Computational Physics 363, 55–78.

Hocevar, M., Sirok, B., Grabec, I., 2004. Experimental turbulent field modeling by visualization and neural networks. Journal of Fluids Engineering 126, 316–322.

Hoffman, J. D., Frankel, S., 2001. Numerical methods for engineers and scientists. CRC Press, NY, USA.

Holmes, P., Lumley, J. L., Berkooz, G., 1998. Turbulence, coherent structures, dynamical systems and symmetry. Cambridge University Press, NY, USA.

Hotelling, H., 1933. Analysis of a complex of statistical variables into principal components. Journal of Educational Psychology 24 (6), 417–441.

Iollo, A., Lanteri, S., Désidéri, J.-A., 2000. Stability properties of POD–Galerkin approximations for the compressible Navier–Stokes equations. Theoretical and Computational Fluid Dynamics 13 (6), 377–396.

Kani, J. N., Elsheikh, A. H., 2017. DR-RNN: A deep residual recurrent neural network for model reduction. arXiv preprint arXiv:1709.00939.

Kerschen, G., Golinval, J.-c., Vakakis, A. F., Bergman, L. A., 2005. The method of proper orthogonal decomposition for dynamical characterization and order reduction of mechanical systems: an overview. Nonlinear Dynamics 41 (1-3), 147–169.

Khibnik, A., Narayanan, S., Jacobson, C., Lust, K., 2000. Analysis of low dimensional dynamics of flow separation. In: Continuation Methods in Fluid Dynamics. Vol. 74. Vieweg, pp. 167–178.

Kim, T.-W., Valdés, J. B., 2003. Nonlinear model for drought forecasting based on a conjunction of wavelet transforms and neural networks. Journal of Hydrologic Engineering 8 (6), 319–328.

Kramer, B., Grover, P., Boufounos, P., Nabi, S., Benosman, M., 2017. Sparse sensing and DMD-based identification of flow regimes and bifurcations in complex flows. SIAM Journal on Applied Dynamical Systems 16 (2), 1164–1196.

Kunisch, K., Volkwein, S., 2002. Galerkin proper orthogonal decomposition methods for a general equation in fluid dynamics. SIAM Journal on Numerical Analysis 40 (2), 492–515.

Kunisch, K., Volkwein, S., 2010. Optimal snapshot location for computing POD basis functions. ESAIM: Mathematical Modelling and Numerical Analysis 44 (3), 569–529.

Lee, C., Kim, J., Babcock, D., Goodman, R., 1997. Application of neural networks to turbulence control for drag reduction. Physics of Fluids 9 (6), 1740–1747.

Lele, S. K., 1992. Compact finite difference schemes with spectral-like resolution. Journal of Computational Physics 103 (1), 16–42.

Levenberg, K., 1944. A method for the solution of certain non-linear problems in least squares. Quarterly of Applied Mathematics 2 (2), 164–168.

Lin, Z., Xiao, D., Fang, F., Pain, C., Navon, I. M., 2017. Non-intrusive reduced order modelling with least squares fitting on a sparse grid. International Journal for Numerical Methods in Fluids 83 (3), 291–306.

Ling, J., Jones, R., Templeton, J., 2016a. Machine learning strategies for systems with invariance properties. Journal of Computational Physics 318, 22–35.

Ling, J., Kurzawski, A., Templeton, J., 2016b. Reynolds averaged turbulence modelling using deep neural networks with embedded invariance. Journal of Fluid Mechanics 807, 155–166.

Ling, J., Templeton, J., 2015. Evaluation of machine learning algorithms for prediction of regions of high Reynolds averaged Navier Stokes uncertainty. Physics of Fluids 27 (8), 085103.

Loève, M., 1955. Probability Theory; Foundations, Random Sequences. New York: D. Van Nostrand Company, NY, USA.

Lorenz, E. N., 1966. Empirical orthogonal functions and statistical weather prediction.

Lu, Y., Zhong, A., Li, Q., Dong, B., 2017. Beyond finite layer neural networks: Bridging deep architectures and numerical differential equations. arXiv preprint arXiv:1710.10121.
Lucia, D. J., Beran, P. S., Silva, W. A., 2004. Reduced-order modeling: new approaches for computational physics. Progress in Aerospace Sciences 40 (1), 51–117.

Lui, H. F., Wolf, W. R., 2019. Construction of reduced order models for fluid flows using deep feedforward neural networks. arXiv preprint arXiv:1903.05206.

Lumley, J., 1967. The structures of inhomogeneous turbulent flow. In: Atmospheric Turbulence and Radio Wave Propagation, A. Yaglom and V. Tatarski, eds. pp. 160–178.

Ly, H. V., Tran, H. T., 2001. Modeling and control of physical processes using proper orthogonal decomposition. Mathematical and Computer Modelling 33 (1-3), 223–236.

Maleewong, M., Sirisup, S., 2011. On-line and off-line POD assisted projective integral for non-linear problems: A case study with Burgers’ equation. International Journal of Mathematical, Computational, Physical, Electrical and Computer Engineering 5 (7), 984–992.

Maulik, R., San, O., 2017. A neural network approach for the blind deconvolution of turbulent flows. Journal of Fluid Mechanics 831, 151–181.

Mignolet, M. P., Przekop, A., Rizzi, S. A., Spottswood, S. M., 2013. A review of indirect/non-intrusive reduced order modeling of nonlinear geometric structures. Journal of Sound and Vibration 332 (10), 2437–2460.

Milton, M., Kountoutsakos, P., 2002. Neural network modeling for near wall turbulent flow. Journal of Computational Physics 182 (1), 1–26.

Mohan, A. T., Gaitonde, D. V., 2018. A deep learning based approach to reduced order modeling for turbulent flow control using LSTM neural networks. arXiv preprint arXiv:1804.09269.

Mohebujaman, M., Rehholz, L. G., Xie, X., Iliescu, T., 2017. Energy balance and mass conservation in reduced order models of fluid flows. Journal of Computational Physics 346, 262–277.

Moosavi, A., Stefanescu, R., Sandu, A., 2015. Efficient construction of local parametric reduced order models using machine learning techniques. arXiv:1511.02909.

Müller, S., Milano, M., Kountoutsakos, P., 1999. Application of machine learning algorithms to flow modeling and optimization. Center for Turbulence Research Annual Research Briefs, Stanford University, Stanford, CA, 169–178.

Narayanan, S., Khibnik, A., Jacobson, C., Kevrekidis, Y., Rico-Martinez, R., Lust, K., 1999. Low-dimensional models for active control of flow separation. In: Proceedings of the 1999 IEEE International Conference on Control Applications. Vol. 2. IEEE, pp. 1151–1156.

Noack, B. R., Morzynski, M., Tadmor, G., 2011. Reduced-order modelling for flow control. Vol. 528. Springer Verlag, Berlin.

Östh, J., Noack, B. R., Krajnović, S., Barros, D., Boréé, J., 2014. On the need for a nonlinear subscale turbulence term in POD models as exemplified for a high-Reynolds-number flow over an Ahmed body. Journal of Fluid Mechanics 747, 518–544.

Peherstorfer, B., Willcox, K., 2016. Data-driven operator inference for nonintrusive projection-based model reduction. Computer Methods in Applied Mechanics and Engineering 306, 196–215.

Ravindran, S. S., 2000. A reduced-order approach for optimal control of fluids using proper orthogonal decomposition. International Journal for Numerical Methods in Fluids 34 (5), 425–448.

Rhinehart, R. R., Su, M., Manimegalai-Sridhar, U., 2012. Leapfrogging and synoptic Leapfrogging: A new optimization approach. Computers & Chemical Engineering 40, 67–81.

Rowley, C. W., Dawson, S. T., 2017. Model reduction for flow analysis and control. Annual Review of Fluid Mechanics 49, 387–417.

Roychowdhury, J., 1999. Reduced-order modeling of time-varying systems. IEEE Transactions on Circuits and Systems II: Analog and Digital Signal Processing 46 (10), 1273–1288.

Sahan, R., Koc-Sahani, N., Albin, D., Liakopoulos, A., 1997. Artificial neural network-based modeling and intelligent control of transitional flows. In: Proceedings of the 1997 IEEE International Conference on Control Applications. IEEE, pp. 359–364.

San, O., Borggaard, J., 2015. Principal interval decomposition framework for POD reduced-order modeling of convective Boussinesq flows. International Journal for Numerical Methods in Fluids 78 (1), 37–62.

San, O., Iliescu, T., 2014. Proper orthogonal decomposition closure models for fluid flows: Burgers equation.
International Journal of Numerical Analysis and Modeling 5, 217–237.
San, O., Iliescu, T., 2015. A stabilized proper orthogonal decomposition reduced-order model for large scale quasigeostrophic ocean circulation. Advances in Computational Mathematics 41 (5), 1289–1319.
San, O., Maulik, R., 2018a. Extreme learning machine for reduced order modeling of turbulent geophysical flows. Physical Review E 97 (4), 042322.
San, O., Maulik, R., 2018b. Machine learning closures for model order reduction of thermal flows. Applied Mathematical Modelling 60, 681–710.
San, O., Maulik, R., 2018c. Neural network closures for nonlinear model order reduction. Advances in Computational Mathematics 44, 1717–1750.
Silveira, L. M., Kamon, M., White, J., 1996. Efficient reduced-order modeling of frequency-dependent coupling inductances associated with 3-D interconnect structures. IEEE Transactions on Components, Packaging, and Manufacturing Technology Part B 19 (2), 283–288.
Singh, A. P., Duraisamy, K., 2016. Using field inversion to quantify functional errors in turbulence closures. Physics of Fluids 28 (4), 045110.
Sirovich, L., 1987. Turbulence and the dynamics of coherent structures. I-Coherent structures. II-Symmetries and transformations. III-Dynamics and scaling. Quarterly of Applied Mathematics 45, 561–571.
Storn, R., Price, K., 1997. Differential evolution—a simple and efficient heuristic for global optimization over continuous spaces. Journal of Global Optimization 11 (4), 341–359.
Taira, K., Brunton, S. L., Dawson, S. T., Rowley, C. W., Colonius, T., McKeon, B. J., Schmidt, O. T., Gordeyev, S., Theofilis, V., Ukeiley, L. S., 2017. Modal analysis of fluid flows: An overview. AIAA Journal 55, 4013–4041.
Tallet, A., Allery, C., Leblond, C., 2016. Optimal flow control using a POD-based reduced-order model. Numerical Heat Transfer, Part B: Fundamentals 70 (1), 1–24.
Tracey, B., Duraisamy, K., Alonso, J., 2013. Application of supervised learning to quantify uncertainties in turbulence and combustion modeling. In: 51st AIAA Aerospace Sciences Meeting including the New Horizons Forum and Aerospace Exposition. p. 259.
Wan, Z. Y., Vlachas, P., Koumoutsakos, P., Sapsis, T., 2018. Data-assisted reduced-order modeling of extreme events in complex dynamical systems. PloS One 13 (5), e0197704.
Wang, C., Liu, J.-G., 2002. Analysis of finite difference schemes for unsteady Navier-Stokes equations in vorticity formulation. Numerical Mathematic 91 (3), 543–576.
Wang, Q., Hesthaven, J. S., Ray, D., 2019. Non-intrusive reduced order modeling of unsteady flows using artificial neural networks with application to a combustion problem. Journal of Computational Physics 384, 289–307.
Wang, Z., Akhtar, I., Borggaard, J., Iliescu, T., 2011. Two-level discretizations of nonlinear closure models for proper orthogonal decomposition. Journal of Computational Physics 230 (1), 126–146.
Wang, Z., Akhtar, I., Borggaard, J., Iliescu, T., 2012. Proper orthogonal decomposition closure models for turbulent flows: a numerical comparison. Computer Methods in Applied Mechanics and Engineering 237, 10–26.
Wang, Z., Xiao, D., Fang, F., Govindran, R., Pain, C. C., Guo, Y., 2018. Model identification of reduced order fluid dynamics systems using deep learning. International Journal for Numerical Methods in Fluids 86 (4), 255–268.
Weller, J., Lombardi, E., Bergmann, M., Iollo, A., 2010. Numerical methods for low-order modeling of fluid flows based on POD. International Journal for Numerical Methods in Fluids 63 (2), 249–268.
Wells, D., Wang, Z., Xie, X., Iliescu, T., 2017. An evolve-then-filter regularized reduced order model for convection-dominated flows. International Journal for Numerical Methods in Fluids 84 (10), 598–615.
Widrow, B., Rumelhart, D. E., Lehr, M. A., 1994. Neural networks: applications in industry, business and science. Communications of the ACM 37 (3), 93–106.
Xiao, D., Fang, F., Buchan, A., Pain, C., Navon, I., Muggeridge, A., 2015a. Non-intrusive reduced order modelling of the Navier–Stokes equations. Computer Methods in Applied Mechanics and Engineering 293, 522–541.
Xiao, D., Fang, F., Pain, C., Hu, G., 2015b. Non-intrusive reduced-order modelling of the Navier–Stokes
equations based on RBF interpolation. International Journal for Numerical Methods in Fluids 79 (11), 580–595.

Xiao, D., Fang, F., Pain, C., Hu, G., 2015c. Non-intrusive reduced-order modelling of the Navier–Stokes equations based on RBF interpolation. International Journal for Numerical Methods in Fluids 79 (11), 580–595.

Xiao, D., Heaney, C., Fang, F., Mottet, L., Hu, R., Bistrian, D., Aristodemou, E., Navon, I., Pain, C., 2019. A domain decomposition non-intrusive reduced order model for turbulent flows. Computers & Fluids 182, 15–27.

Xie, X., Mohebujjaman, M., Rebholz, L., Iliescu, T., 2018a. Data-driven filtered reduced order modeling of fluid flows. SIAM Journal on Scientific Computing 40 (3), B834–B857.

Xie, X., Wells, D., Wang, Z., Iliescu, T., 2017. Approximate deconvolution reduced order modeling. Computer Methods in Applied Mechanics and Engineering 313, 512–534.

Xie, X., Wells, D., Wang, Z., Iliescu, T., 2018b. Numerical analysis of the Leray reduced order model. Journal of Computational and Applied Mathematics 328, 12–29.

Xu, K., Darve, E., 2019. The neural network approach to inverse problems in differential equations. arXiv preprint arXiv:1901.07758.

Zhang, G., Patuwo, B. E., Hu, M. Y., 1998. Forecasting with artificial neural networks: The state of the art. International Journal of Forecasting 14 (1), 35–62.