VALIDATION OF SPALLATION MODELS:
AN APPROACH*

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To understand and describe the proton induced spallation reactions, a large number of computer codes have been proposed. Various quantitative tests are used in literature to judge the agreement between model calculations and experimental data. The judgement is based on the magnitude of the deviation of the tests from their expected values in the case of the perfect agreement. However, the expected values of the tests and their standard deviations are usually not well known. Thus the conclusions may be ambiguous. It is proposed to calculate the expected values and standard deviations of the tests by Monte Carlo method and to use the tests in their standardized form.

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1. Introduction

Proton induced spallation reactions cover a wide range of applications in various branches of physics, e.g., particle transport codes, accelerator based nuclear waste transmutation, neutron spallation targets, astrophysics etc. It becomes very important to find a theoretical model which can describe a complete set of reaction outcomes. The tests which are most frequently used in literature to judge the quality of data description by individual models are listed in Table I.

The expectation value of each test corresponds to the perfect agreement of model $\sigma^{\text{calc}}$ and experimental $\sigma^{\text{exp}}$ cross sections. However, it should be taken into consideration that the experimental data are biased by statistical errors, thus they are statistical variables. Therefore, the perfect agreement should be meant as

$$E(\sigma^{\text{exp}}) = \sigma^{\text{calc}}.$$  

(1.1)

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The most frequently used tests for validation of spallation models. The symbols $\sigma_{i}^{\text{exp}}$ and $\Delta\sigma_{i}^{\text{exp}}$ correspond to experimental cross sections and their statistical errors, whereas $\sigma_{i}^{\text{calc}}$ denotes model cross sections.

| Formula of the test | Expectation value | Ref. (equation) |
|---------------------|-------------------|-----------------|
| $H = \left[ \frac{1}{N} \sum_{i=1}^{N} \left( \frac{\sigma_{i}^{\text{exp}} - \sigma_{i}^{\text{cal}}}{\Delta\sigma_{i}^{\text{exp}}} \right)^{2} \right]^{1/2}$ | $E(H) \approx 1$ | [1] (1) |
| $D = \frac{1}{N} \sum_{i=1}^{N} \left| \frac{\sigma_{i}^{\text{exp}} - \sigma_{i}^{\text{cal}}}{\sigma_{i}^{\text{cal}}} \right|$ | $E(D) \approx 0$ | [1] (3) |
| $R = \frac{1}{N} \sum_{i=1}^{N} \frac{\sigma_{i}^{\text{exp}}}{\sigma_{i}^{\text{cal}}}$ | $E(R) \approx 1$ | [1] (5) |
| $L = \left[ \sum_{i=1}^{N} \left( \frac{\sigma_{i}^{\text{calc}}}{\Delta\sigma_{i}^{\text{exp}}} \right)^{2} \left( \frac{1}{\sigma_{i}^{\text{calc}}} \right)^{2} \left/ \sum_{i=1}^{N} \left( \frac{\sigma_{i}^{\text{calc}}}{\Delta\sigma_{i}^{\text{exp}}} \right)^{2} \right]^{1/2} \right.$ | $E(L) \approx 0$ | [1] (8) |

The expectation values of tests listed in the table are estimated assuming that the following equation holds

$$E\left(\text{Test}\left(\sigma^{\text{exp}}\right)\right) = \text{Test}\left(E\left(\sigma^{\text{exp}}\right)\right)$$

which is strictly valid only for linear functions but actually serves as a reasonable approximation. The ranking of models is usually done by a comparison of the magnitude of the deviation of the test value from its value expected for the case of the perfect agreement with data. Such an approach allows to determine the relative success of different model descriptions but cannot answer the question whether the description is accurate enough to be treated as a proof of the model adequacy. This is caused by the lack of information on the standard deviations of most of the tests.

2. The proposed approach

In the present work, it is proposed: (i) to evaluate the expectation value of each test $E(\text{Test})$ and its standard deviation $\sigma(\text{Test})$ by the Monte Carlo method for the case of the perfect agreement of the model and data cross sections (defined by Eq. (1.1)), and (ii) to use instead the statistical tests
listed in Table I, the standardized test evaluated according to the formula

\[ z(\text{Test}) \equiv \frac{\text{Test} - E(\text{Test})}{\sigma(\text{Test})}. \tag{2.1} \]

The Monte Carlo sampling of the test values is performed under two assumptions concerning experimental data:

(i) The experimental data are independent Gaussian variables with standard deviation equal to statistical error of data. This assumption is fulfilled when the statistics of experimental events is large enough. When the statistics is poor, one has to use the Poisson distribution.

(ii) Expectation value of the experimental cross section is equal to the model cross section (Eq. (1.1)). The alternative possibility is to use the experimental cross section as the estimate of its expectation value.

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**Fig. 1.** The open dots depict histograms of H-test calculated for two energy ranges of the experimental spectrum of deuterons from Al\((p,d)\) reaction \((E_d < 30 \text{ MeV} \text{ and } E_d > 30 \text{ MeV})\) measured at \(\vartheta = 100^\circ\) for proton beam energy 1.2 GeV [2]. The lines represent the analytical probability distributions of the H-test Eq. (2.2). The 6 data points are lying in the low energy region, whereas 58 data points in the high energy region. The expectation values and standard deviations estimated from the histograms and evaluated using Eqs. (2.3)–(2.4) are also listed.
To check reliability of this method, the analytically derived probability distribution, expectation value and standard deviation of the H-test were compared with the quantities obtained by the Monte Carlo sampling. It may be shown that with the above assumptions the probability distribution of the H-test and its parameters have the following form:

\[
h(H) = \frac{N^{N/2}}{2^{(N/2)-1} \Gamma(N/2)} \exp\left(\frac{-NH^2}{2}\right) H^{N-1},
\]

(2.2)

\[
E(H) = \sqrt{\frac{2}{N}} \frac{\Gamma\left(\frac{N+1}{2}\right)}{\Gamma\left(\frac{N}{2}\right)},
\]

(2.3)

\[
\sigma(H) = \sqrt{1 - E(H)^2}.
\]

(2.4)

The comparison of histograms obtained with the Monte Carlo method and analytical formulae of the H-test are presented in Fig. 1. As can be seen, the agreement between analytic distributions and histograms is perfect.

3. Results

As an example of the proposed method, the standardized tests were calculated for the experimental spectrum of the Al($p$,d) reaction at beam energy 1.2 GeV measured at $\vartheta = 100^\circ$ [2] with results of calculations performed by the means of INCL4.6 intranuclear cascade code coupled to SMM code. The values have been calculated separately for $E_d < 30$ MeV and $E_d > 30$ MeV because in the first energy region evaporation dominates, whereas for the second region the prompt emission from the cascade stage is prevailing (due to coalescence of escaping neutron and proton). The following standardized values of the test have been obtained for $E_d < 30$ MeV: 7.01, 8.52, 6.25, 7.73 for H, D, R, and L tests, respectively. The corresponding values for $E_d > 30$ MeV are equal to: 28.48, 7.59, $-7.03$, and 23.69. For $E_d < 30$ MeV, where the relative errors of all data points are almost the same, all standardized tests have nearly the same values. For $E_d > 30$ MeV, the relative errors of the data vary by one order of magnitude what seems to be a reason of different values of standardized test H and L in respect to R and D.

It is worth emphasizing that values of nonstandardized tests are quite different even for $E_d < 30$ MeV: 2.907, 0.109, 1.097, 0.113 for H, D, R, and L tests, respectively. This does not allow to judge whether they contain the compatible information.
It is reasonable to conjecture that the probability distributions of standardized tests are close to the standard normal distribution $N(0,1)$ because of the central limit theorem. Therefore, standardized tests may be easily used for making the decision on acceptance of the hypothesis claiming the good agreement of the model with the data. In the present example, this hypothesis may be rejected on the significance level smaller than 0.0027.

The proposed approach allows to make objective, quantitative comparisons concerning the quality of data reproduction by different theoretical spallation models which is not possible using usual, non-standardized tests.

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