Electron-Deficient Zn-N$_6$ Configuration Enabling Polymeric Carbon Nitride for Visible-Light Photocatalytic Overall Water Splitting

Daming Zhao$^{1,2}$, Yiqing Wang$^1$, Chung-Li Dong$^3$, Fanqi Meng$^4$, Yu-Cheng Huang$^3$, Qinghua Zhang$^4$, Lin Gu$^4$, Lan Liu$^1$, Shaohua Shen$^1$

HIGHLIGHTS

- Atomically dispersed Zn-anchored 3D sponge-like polymeric carbon nitride (Zn-PCN) characteristic of a unique Zn-N$_6$ electron-deficient configuration is synthesized via an intermediate coordination strategy.
- The electron-deficient Zn-N$_6$ configuration contributes to enhanced electron excitation, accelerated charge separation and transfer as well as reduced overpotentials of water redox reactions.
- The obtained Zn-PCN realizes photocatalytic overall water splitting to stoichiometrically produce H$_2$ and O$_2$ with good durability under visible light.

ABSTRACT Despite of suitable band structures for harvesting solar light and driving water redox reactions, polymeric carbon nitride (PCN) has suffered from poor charge transfer ability and sluggish surface reaction kinetics, which limit its photocatalytic activity for water splitting. Herein, atomically dispersed Zn-coordinated three-dimensional (3D) sponge-like PCN (Zn-PCN) is synthesized through a novel intermediate coordination strategy. Advanced characterizations and theoretical calculations well evidence that Zn single atoms are coordinated and stabilized on PCN in the form of Zn-N$_6$ configuration featured with an electron-deficient state. Such an electronic configuration has been demonstrated contributive to promoted electron excitation, accelerated charge separation and transfer as well as reduced water redox barriers. Further benefited from the abundant surface active sites derived from the 3D porous structure, Zn-PCN realizes visible-light photocatalysis for overall water splitting with H$_2$ and O$_2$ simultaneously evolved at a stoichiometric ratio of 2:1. This work brings new insights into the design of novel single-atom photocatalysts by deepening the understanding of electronic configurations and reactive sites favorable to excellent photocatalysis for water splitting and related solar energy conversion reactions.

KEYWORDS Zn single atoms; Polymeric carbon nitride; Overall water splitting; Photocatalysis

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1 Introduction

Spontaneously solar-driven water splitting into $H_2$ and $O_2$ via artificial photosynthesis is very appealing for addressing the dilemma from fossil fuels and environmental issues [1, 2]. Extensive efforts have been made in the past few decades to develop various semiconducting materials as photocatalysts, mainly including metal oxides (e.g., $TiO_2$, $Fe_2O_3$, $WO_3$, $BiVO_4$, $SrTiO_3$), metal sulfides (e.g., $CdS$, $Cd_1_xZn_1_xS$, $CuInS_2$, $CuZnSnS_4$), and metal (oxy)nitrides (e.g., $Ta_3N_5$, $TaON$, $GaN/ZnO$, $LaTiO_2N$) [3, 4]. Recently, polymeric semiconductors with adjustable molecular structures and tunable band structures have been documented as the promising alternatives to inorganic photocatalysts for water splitting [5–7]. Typically, as an organic semiconductor mainly composed of carbon and nitrogen, polymeric carbon nitride (PCN) with a two-dimensional (2D) conjugated structure has inspired extensive research in photocatalysis for $H_2$ and/or $O_2$ production, owing to its abundant precursors, environmental friendliness, chemical stability and proper band potentials for water redox reactions [8–11], since its pioneering demonstration by Wang et al. [12]. However, the intrinsic drawbacks, especially the high charge-carrier recombination rate and the sluggish surface reaction kinetics, largely restrict the photocatalytic efficiency of PCN [8].

To conquer these limitations, a variety of strategies including heteroatoms doping, defect engineering, heterojunction constructing, morphology design, etc., [13–18] have emerged to modify PCN for photocatalytic water splitting, with performances continually hitting record highs.

As a hot spot of the current research in catalysis, metal single-atom catalysts with nearly 100% atomic utilization and superb catalytic properties have received increasing attentions in the fields of thermocatalysis, electrocatalysis and photocatalysis [19–23]. Given the large surface free energy induced by the extremely small size, a support that could chemically interact and then highly disperse single atoms is always desired to prohibit the aggregation and ensure the stability of single-atom catalysts. With abundant electron-rich pyridinic nitrogen atoms in the tri-s-triazine units, PCN has been believed as an ideal support for stabilizing single atoms, by providing sufficient coordination sites to anchor metal atoms via the formed strong covalent metal-N bonds [9, 24]. In turn, the physicochemical property of PCN support would be concomitantly changed with the aid of the coordinated single atoms. Thanks to the fine-tunable electronic structure of PCN and the high catalytic activity of single-atom active sites, PCN-based metal single-atom photocatalysts have shown great potentials in photocatalytic water splitting, with charge transfer processes promoted and surface reaction kinetics accelerated [25–29]. For instance, Guo group developed a $PH_3$-assisted thermolysis strategy to create a phosphorus-coordinated noble metal single-atom coordination on PCN [25]. With an electron-rich feature in the $MP_2$ ($M=Ru, Rh, Pd$) coordination contributing to reduced reaction barriers, the obtained $MP_2/PCN$ exhibited excellent visible-light photocatalytic $H_2$ production activity and stability. Bi and co-workers achieved a significantly improved $H_2$ production activity by anchoring single-atom Pt onto PCN via photochemical deposition, and elucidated the separated reduction and oxidation reactions sites by discovering the dynamic variations of Pt–N bond/Pt$^0$ species and the corresponding C-N bond/C = N bond in single-atom Pt/PCN during the photocatalytic water splitting process using synchronous illumination X-ray photoelectron spectroscopy [26]. Shi et al. introduced Ni single-atom sites onto PCN via a freezing-deposition method and precisely modulated the oxidation state of Ni into the intermediate state with a $Ni^{2+}$/Ni$^0$ ratio of 2, which brought abundant unpaired $d$-electrons in the Ni/PCN photocatalyst [27]. Such an electronic configuration resulted in promoted light response, conductivity, charge separation and mobility, thus enhancing the photocatalytic water reduction performance. With significant advances achieved for PCN-based metal single-atom photocatalysts, the insightful understanding of structure–activity relationship is still essential for the development and design of novel single-atom photocatalysts toward efficient solar-hydrogen conversion.

Herein, atomically dispersed Zn-coordinated three-dimensional (3D) sponge-like PCN (Zn-PCN) was synthesized via a novel intermediate coordination strategy, with single-atom loading content up to 4.79% and specific surface area reaching 224.99 $m^2$ g$^{-1}$. In comparison with bulk PCN (BCN), the obtained Zn-PCN achieves significantly improved photocatalytic $H_2$ production activity from water splitting in the presence of sacrificial agent. Surprisingly, Zn-PCN is capable of photocatalytic overall water splitting under visible light, producing $H_2$ and $O_2$ at a stoichiometric ratio of 2:1, with $H_2$ production rate reaching 35.2 $\mu mol$ h$^{-1}$ g$^{-1}$. Along with the 3D porous structure beneficial to abundant surface active sites, experimental
and calculation results demonstrate that Zn single atoms are steadily anchored on PCN via a unique Zn-N₆ electron-deficient coordination, which induces a midgap energy level, increases the carrier concentration, and strengthens the charge polarization, contributing to enhanced electron excitation, accelerated charge separation and transfer as well as reduced overpotentials of water redox reactions. This work provides a new idea for the design of novel single-atom photocatalysts toward solar water splitting with in-depth understanding of the structure–activity relationship from the level of atomic coordination and electronic configuration.

2 Experimental Section

2.1 Materials

Melamine (C₃H₆N₆, ≥ 99.0%), hydrochloric acid (HCl, ~ 36.0–38.0 wt%), sodium sulfate (Na₂SO₄, ≥ 99.0%), and triethanolamine (TEOA, analytical reagent) were purchased from Sinopharm Chemical Reagent Co., Ltd. Zinc carbonate (ZnCO₃, 97%) and chloroplatinic acid hexahydrate (H₂PtCl₆·6H₂O, ACS reagent) were purchased from Alfa Aesar. Heavy-oxygen water (H₂¹⁸O, 99%) was purchased from 3A chemicals Co., Ltd. Ultrahigh-purity argon (Ar, ≥ 99.999%) was provided by Shaanxi Xinkang Medical Oxygen Co., Ltd. All the materials were used as received without further purification. Deionized water, with a resistivity of 18.25 MΩ cm, was used throughout the experiments.

2.2 Synthesis of Bulk PCN

Bulk PCN (BCN) was synthesized by thermal polymerization of melamine. Typically, 2 g of melamine was put into a covered alumina crucible and calcined at 520 °C in Ar for 4 h with a ramping rate of 5 °C min⁻¹. After naturally cooled to room temperature, the yellow BCN was collected and ground into powder.

2.3 Synthesis of Atomically Dispersed Zn-PCN

Atomically dispersed Zn-coordinated 3D sponge-like PCN was synthesized through an intermediate coordination method. In detail, 2 g of melamine was mixed with a certain amount of ZnCO₃ (1, 2, 4 g) and then calcined at 400 °C in Ar for 2 h with a ramping rate of 5 °C min⁻¹. After naturally cooled to room temperature, the resultant solid mixtures were added into excessive HCl aqueous solution (1.0 mol L⁻¹) and vigorously stirred for 12 h, followed by washing with deionized water until the supernatant was neutral. After freeze drying (24 h), the obtained powders were calcined at 520 °C in Ar for 4 h with a ramping rate of 5 °C min⁻¹. The final products were denoted as Zn-PCN(x) (x represents the weight contents of Zn single atoms, x = 2.09%, 3.63%, 4.79%).

2.4 Characterization

The weight contents of zinc were measured by an inductively coupled plasma mass spectrometer (NexION 350D, PerkinElmer). Transmission electron microscopy (TEM) images were recorded on a FEI Tecnai G2 F30 S-Twin transmission electron microscope at an accelerating voltage of 300 kV. An OXFORDMAX-80 energy-dispersive X-ray detector, mounted in this microscope, was used to conduct the elemental analysis. The high-angle annular dark-field scanning transmission electron microscopy (HAADF-STEM) was performed on a JEOL JEM-ARM200F transmission electron microscope operated at 200 kV and equipped with double spherical aberration correctors. Scanning electron microscopy (SEM) images were obtained by a JEOL JSM-7800F field-emission scanning electron microscope. N₂ adsorption–desorption isotherms were conducted using a surface area and porosity analyzer (ASAP 2020, Micromeritics) at 77 K after degassing the samples at 150 °C for 4 h. The specific surface area and pore volume were determined by the Brunauer–Emmett–Teller method. The pore size distribution was calculated using the Barrett–Joyner–Halenda method. X-ray diffraction (XRD) patterns were collected on a PANalytical X’pert Pro MPD diffractometer operated at 40 kV and 40 mA using Ni-filtered Cu Kα irradiation (wavelength = 1.5406 Å). Fourier transform infrared (FTIR) spectra were recorded using a Bruker Vertex 70 FTIR spectrometer using the KBr pellet technique. The amount of sample used for each FTIR spectrum was kept at around 1 mg. X-ray photoelectron spectroscopy (XPS) data were obtained on a Kratos Axis-Ultra DLD.
spectrometer with a monochromatized Al Kα line source (150 W). All the binding energies were referenced to the C 1s peak at 284.8 eV. Ultraviolet–visible–near infrared diffuse reflectance spectra were recorded on a Cary Series ultraviolet–visible–near infrared spectrophotometer (Agilent Technologies). The surface photovoltage (SPV) spectra were obtained on a self-made measurement system, equipped with a 500 W Xe lamp (CHF-XM500, Beijing PerfectLight), a grating monochromator (Omni-λ3007, Zolix), a lock-in amplifier (SR830-DSP) with an optical chopper (SR540) and a photovoltaic cell. The steady-state photoluminescence (PL) emission spectra and time-resolved transient PL decay spectra were measured at room temperature using a PTI QM-4 fluorescence spectrophotometer. Considering the existence of three radiative lifetimes, the PL decay curve of PCN was triexponentially fitted using equation [16]:

$$I(t) = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} + A_3 e^{-t/\tau_3}$$

(1)

$$\tau_{avg} = \frac{(A_1 \tau_1^2 + A_2 \tau_2^2 + A_3 \tau_3^2)}{(A_1 \tau_1 + A_2 \tau_2 + A_3 \tau_3)}$$

(2)

where $A_1$, $A_2$, and $A_3$ represent the normalized amplitudes of each decay component, respectively, $\tau_1$, $\tau_2$, and $\tau_3$ are values of the lifetime components, respectively, $\tau_{avg}$ is average lifetime of photogenerated charge carriers.

2.5 Synchrotron-Based X-ray Spectroscopic Measurements and Data Analysis

The synchrotron-based X-ray spectroscopic measurements at the C, N and Zn K-edge were carried out at beamline BL20A, at the National Synchrotron Radiation Research Center. The acquired EXAFS data were processed according to the standard procedures using the ATHENA module implemented in the IFEFFIT software packages [25]. The $k^3$-weighted $\chi(k)$ data in the $k$-space ranging from 2.893 to 10.99 Å⁻¹ were Fourier transformed to real ($R$) space using Hanning windows (dk = 1.0 Å⁻¹) to separate the EXAFS contributions from different coordination shells. To obtain the detailed structural parameters around Zn atom in the as-prepared samples, the quantitative curve fittings were performed for the Fourier transformed $k^3\chi(k)$ in the $R$-space using the ARTEMIS module of IFEFFIT. Effective back-scattering amplitudes $F(k)$ and phase shifts $\Phi(k)$ of all fitting paths were calculated by the ab initio code FEFF8.0. During the fitting analysis, the amplitude reduction factor $S_0^2$ was fixed to the best-fit value of 0.8011, which was determined from fitting the reference sample of Zn foil. To fit the data of Zn-PCN, the Zn-N coordination number, interatomic distance ($R$), Debye–Waller factor ($\sigma^2$) and inner potential shift ($\Delta E_0$) were allowed to vary. The obtained structural parameters are summarized in Table S2.

2.6 Density Functional Theory Calculations

First-principles density functional theory (DFT) calculations were performed using the CASTEP code [30]. The exchange–correlation effects were treated in generalized gradient approximation (GGA) with the Perdew–Burke–Ernzerhof (PBE) potential. The interaction between ionic core and valence electrons was simulated by the ultrasoft pseudopotentials. The kinetic energy cut-off was chosen to be 500 eV. Brillouin zone integration was sampled with $3 \times 3 \times 3$ and $3 \times 3 \times 1$ Monkhorst–Pack mesh K-points for BCN and Zn-PCN calculation, respectively. The slab thickness for three-layer structure models is 6.54 Å. The slabs were set by a vacuum region of 20 Å. The adsorbates and the top layers of the surfaces are allowed to fully relax and optimize until convergence to $10^{-5}$ eV in total energy and 0.01 eV Å⁻¹ in the forces.

2.7 Photoelectrochemical Measurements

The photoelectrochemical properties were measured on an electrochemical workstation (CHI760D) in a three-electrode cell system under light irradiation, in which platinum foil and Ag/AgCl electrode were used as the counter and reference electrodes, respectively. The as-prepared photocatalysts were electrodeposited on fluoride tin oxide (FTO) glass substrates and used as the working electrodes. In detail, 40 mg of photocatalysts were ultrasonically dispersed in 50-mL acetone solution of I$_2$ (10 mg). One piece of clean FTO glass (1 × 1.5 cm$^2$) and one platinum foil in parallel, immersed in the above dispersion, were applied with a potential of −5 V for 5 min. The immersed area of the FTO glass was fixed at 1 × 1 cm$^2$. Then, the deposited electrode was dried at 120 °C for 30 min to remove the I$_2$ residue. 0.5 M Ar-saturated Na$_2$SO$_4$ aqueous solution was used as the electrolyte. A 500 W Xe lamp
(CHF-XM500, Beijing PerfectLight) with adjustable power settings through an air mass 1.5 global (AM1.5) filter was used as light source (100 mW cm⁻²). The illumination area of working electrode was fixed at 0.785 cm². The transient photocurrent response with an interval of 10 s on/off switch was measured at an applied potential of 0.2 V versus Ag/AgCl. Electrochemical impedance spectra (EIS) were recorded at an applied potential of 0.8 V versus Ag/AgCl. The frequency varied from 10 kHz to 0.1 Hz, and the amplitude of the sinusoidal wave was set at 5 mV. The Mott–Schottky plots were measured with 10 mV amplitude at different frequencies (1000, 2000 and 3000 Hz).

2.8 Photocatalytic Measurements

Photocatalytic H₂ production reactions were performed in a hermetic Pyrex reactor. In a typical reaction, 10 mg of photocatalysts were dispersed in 100-mL aqueous solution of 10 vol% TEOA using a magnetic stirrer with constant rotational velocity. 1 wt% Pt as cocatalyst was photodeposited on the photocatalyst from the precursor of H₂PtCl₆·6H₂O. Ar was purged through the reactor for 30 min before reaction to remove the residual air. A 300 W Xe lamp (PLS-SXE300, Beijing PerfectLight) with a 420-nm cut-off filter (λ > 420 nm) was used as light source to trigger the photocatalytic H₂ generation. The temperature of the reaction solution was kept at 35 °C via a circulating water pump during the whole experiment. Evolved H₂ was measured through a gas chromatograph equipped with a thermal conduction detector (5 Å molecular sieve column, Ar as carrier gas). Photocatalytic overall water splitting reactions were conducted under the same photocatalytic reaction condition without any carrier scavenger. Blank experiments revealed no appreciable gas production without irradiation or photocatalysts. Photocatalytic H₂¹⁸O (99%) splitting measurements were detected on an Agilent 7890A-5975C gas chromatograph–mass spectrometer.

3 Results and Discussion

3.1 Catalyst Synthesis and Morphology

Zn-PCN was synthesized through an intermediate coordination method by using melamine as the precursor of PCN, and ZnCO₃ as zinc sources and also the template for 3D sponge-like structure (see Experimental Section for details). As shown in Fig. 1a, through the first-step calcination, a mixture of melem monomer and ZnO was produced from melamine and ZnCO₃. With ZnO leached by hydrochloric acid (HCl) solution, the dried mixture, in which residual Zn²⁺ species might coordinate with the pyridinic nitrogen atoms of melem monomer [9, 31], was then calcined again to obtain the Zn-PCN products via the polymerization of melem intermediates with the coordinated Zn single atoms. By changing the amounts of ZnCO₃ precursor, the weight contents (x) of Zn single atoms in Zn-PCN(x) could be readily tuned (x = 2.09%, 3.63%, 4.79%), as quantified by inductively coupled plasma mass spectroscopy. As compared to BCN synthesized by direct thermal polymerization of melamine (Fig. S1a-b), the obtained Zn-PCN presents a unique 3D sponge-like morphology (Figs. S1c-d), with the produced ZnO species acting as hard templates and the released NH₃ and CO₂ gases acting as soft templates during the first-step calcination. As learned from the N₂ adsorption–desorption measurements (Fig. S2 and Table S1), Zn-PCN owns a typical mesoporous structure with average pore diameter of ~30 nm. Depending on the increasing amounts of ZnCO₃ precursor (i.e., Zn loading contents), the specific surface area gradually increases, reaching 224.99 m² g⁻¹ as high for Zn-PCN(4.79%), which is almost 24 times that for BCN (9.46 m² g⁻¹). Such a well-developed 3D porous structure with increased specific surface area and inhibited agglomeration could be confirmed by the weakened XRD diffraction peaks (Fig. S3) [15, 32], which would provide abundant surface active sites for accelerating photocatalytic reaction [33, 34]. With nanoparticle and clusters of ZnO or other Zn-based compounds hardly identified in TEM image (Fig. 1b) and XRD patterns (Fig. S3), the high dispersion of Zn single atoms on PCN could be further evidenced from the spherical aberration-corrected high-angle annular dark-field scanning transmission electron microscopy (HAADF-STEM) images (Fig. 1c, d). One could easily discover densely and evenly isolated bright spots with sizes <0.2 nm (highlighted by yellow circles) that could be assigned to the Zn atoms heavier than the non-metal atoms such as N and C in PCN [25–29], which unambiguously confirms the good dispersion of Zn atoms with high density in Zn-PCN. With the existence of Zn element further affirmed by energy-dispersive X-ray spectroscopy (EDS) spectrum, elemental mappings again reveal the homogeneous dispersion of Zn atoms along with elemental C and N in Zn-PCN (Fig. 1e,
These microscopy observations provide solid evidences for the high dispersion of Zn single atoms anchored on the PCN support, with the molecular structural characteristics of both tri-s-triazine ring and trigonal C-N(-C)-C/bridging C-NH-C units well-maintained during the intermediate coordination and polymerization process for the introduction of Zn single atoms (Fig. S4).

### 3.2 Structure Determination

The chemical state and coordination environment of Zn single atoms in Zn-PCN were determined by synchrotron-based X-ray spectroscopic measurements collected at the Zn K-edge of Zn-PCN and some references (ZnO and Zn foil). As shown in the Zn K-edge X-ray absorption near-edge structure (XANES) spectra (Fig. 2a), the absorption threshold for Zn-PCN is significantly higher than that for Zn foil and even higher than that for ZnO, indicating that Zn single atoms in Zn-PCN exist at an oxidation state higher than that of Zn$^{2+}$ and hold more unoccupied states than Zn$^{2+}$ [29, 31], as further confirmed by that the characteristic peaks of Zn 2p$^{3/2}$ and Zn 2p$^{1/2}$ levels for Zn-PCN shift to higher binding energies as compared to ZnO in the Zn 2p XPS spectra (Fig. 2b) [35]. Considering that Zn$^{2+}$ has a 3$d^{10}$ electron configuration, Zn
single atoms should possess a partially vacant 3d orbital (inset in Fig. 2a), attributed to the electron transfer from Zn single atoms to PCN support. Note the two O 1s XPS signals at 530.2 and 531.8 eV for ZnO as reference (Fig. S5), assigned to the lattice oxygen in wurtzite structure and defective oxygen, respectively [36]. Zn-PCN only displays a weak O 1s XPS signal assigned to absorbed oxygen (532.1 eV) [37], excluding the Zn–O interaction in Zn-PCN. In the Fourier transform extended X-ray absorption fine structure (FT-EXAFS) spectra (Fig. 2c),
without signals detected for Zn–O interaction (~ 2.9 Å) or Zn-Zn interaction (~ 2.3 Å), only a distinct peak attributed to Zn-N interaction (~ 1.5 Å) could be observed for Zn-PCN [38, 39], suggesting that Zn single atoms should be coordinated with the unsaturated N sites in the tri-s-triazine units of Zn-PCN with an electron-deficient state. As further revealed by the wavelet transforms (WT) of EXAFS (Fig. 2d), rather than the peaks at \( k = \sim 8.0 \text{ Å}^{-1} \) observed for ZnO and at \( k = \sim 7.0 \text{ Å}^{-1} \) for metallic Zn, Zn-PCN presents a single peak at \( k = \sim 5.5 \text{ Å}^{-1} \) in the contour plot, which again corroborates the Zn-N coordination in Zn-PCN. With the quantitative structural parameters of Zn single atoms in Zn-PCN confirmed by the fitted FT-EXAFS curve, it is convincing that Zn coordinated by six pyridinic N (i.e., Zn-N_6) should be the dominant coordination structure in Zn-PCN (Fig. 2e and Table S2).

In addition to the characterization of Zn single atoms, XANES measurements at C K-edge and N K-edge were also performed to further elucidate the local atomic and electronic structure changes in Zn-PCN versus BCN. As shown in Fig. 2f, the C K-edge peak at 287.4 eV is associated with the electron transition from C 1s to 2p \( \pi^* \) orbital related to N–C = N \( \text{sp}^2 \) hybridized states in BCN; the N K-edge peaks at 399.1 and 401.8 eV correspond to the N1 site of C-N = C coordination structure in a tri-s-triazine unit and the N2 site of N-3C bridging structure connecting with three tri-s-triazine units, respectively [32]. One could observe that both the C K-edge and N K-edge peaks of Zn-PCN shift toward photon energy with lower peak intensity as compared to those of BCN (Fig. 2f), meaning the decreased \( \pi^* \) characteristics of Zn-PCN due to the chemical coordination between tri-s-triazine units and the anchored Zn single atoms [16], as evidenced in the Zn K-edge XANES spectra (Fig. 2a). Thereafter, the less unoccupied 2p states at tri-s-triazine units of Zn-PCN could generally accelerate the electron transfer from Zn single atoms to tri-s-triazine units and thus benefits photocatalysis [31].

### 3.3 Density Functional Theory Calculations

DFT calculations were performed to better understand the electronic configuration evolution of PCN induced by the anchored Zn single atoms. With monolayer and three-layer structure models set for 3D sponge-like Zn-PCN and BCN, respectively, the band structure and density of states (DOS) calculation results demonstrate that both C 2p and N 2p orbitals contribute to the conduction band (CB), while the valence band (VB) is dominated by N 2p orbitals for PCN (Fig. 3a-c), which is well-consistent with previous reports [40]. Upon the introduction of Zn single atoms, Zn 3d orbitals greatly contribute to the CB of Zn-PCN, with a new Zn-related midgap energy level appearing at ~ 2.07 eV above the VB (Fig. 3d–f), coinciding with the extended absorption tails seen in the ultraviolet–visible-near infrared diffuse reflectance spectra (UV–vis-NIR DRS, Fig. S6a) as well as the experimentally determined band structures (Fig. S6b–d). Meanwhile, the bandgap of Zn-PCN is widened as compared to BCN (Fig. 3b, e), agreeing with the blueshift in optical absorption edge upon the introduction of Zn single atoms (Fig. S6a), which should be related to the well-developed 3D porous structure inhibiting the agglomeration of Zn-PCN [32]. It should be noted that the different Zn single-atom loading contents do not obviously change the bandgap of Zn-PCN (Fig. S7a-f). However, the high Zn single-atom loading content would give rise to a small work function (\( \Phi \)) of Zn-PCN (Fig. S7g), with the work function calculated for Zn-PCN (4.25 eV) much smaller than that of BCN (4.66 eV) (Fig. 3g), which suggests the higher carrier concentration, contributing to the better charge transfer capability and the stronger electron excitation in Zn-PCN (Fig. S7h–i) [27, 41]. The differential charge density map (Fig. 3h) witnesses the charge redistribution in Zn-PCN with holes mainly enriched at Zn single atoms and electrons accumulated at the adjacent C and N atoms, further implying the electron-deficient feature of Zn single atoms. Such localized charge distribution would spatially separate the redox active sites from each other, inhibiting the recombination of photogenerated electron–hole pairs and the back reaction of \( \text{H}_2 \) and \( \text{O}_2 \) evolved at surface [31].

### 3.4 Electron Excitation and Charge Transfer Properties

Above calculation results indicate that the introduced Zn single atoms could benefit the electron excitation by modulating the electronic configuration in Zn-PCN, which was further explored by the C K-edge and N K-edge XANES spectra recorded both in dark and under illumination to monitor
the photoexcited electron transition in Zn-PCN. As previously reported [42], PCN is composed of tri-s-triazine units interconnected with tertiary nitrogen atoms, where C-N $sp^3$ hybridization constitutes the high-energy $\sigma$ and $\sigma^*$ molecular orbitals, while C-N $sp^2$ hybridization gives rise to the lower-energy $\pi$ bonding and $\pi^*$ antibonding orbitals. In addition, the unpaired electrons on pyridinic N atoms create a lone pairs (LP) orbital above the $\pi$ bonding orbital, and vacancies or dopants would induce a midgap energy level. These orbitals form the occupied and unoccupied states, respectively (Fig. 4a). It is clear that with the introduction of Zn single atoms, the irradiation-dark intensity variations in both C K-edge and N K-edge spectra are more obvious in Zn-PCN than in BCN (Fig. 4b), supporting the promoted electron excitation from LP states to $\pi^*$ states in the CB of Zn-PCN [43], which should be related to the increased carrier concentration caused by the electron injection from Zn single atoms to PCN [27]. Then, the charge transfer processes in Zn-PCN was investigated by photoluminescence (PL) spectroscopy. As shown in Fig. 4c, a strong PL emission peak centered at ~465 nm could be observed for BCN, which is related to the electron transition between $\pi^*$ antibonding orbital and LP orbital [12, 42]. By comparison, Zn-PCN show a dramatical PL quenching (Fig. 4c), implying that the radiative recombination of photogenerated electron–hole pairs in Zn-PCN could be substantially inhibited by the introduced midgap energy level (Fig. 4a) and the strengthened charge polarization (Fig. 3h) by the anchored Zn single atoms [25, 26]. Note that the gradual blueshift in PL emission peak depends on the increasing amounts of Zn single atoms, inferring the widened bandgaps in accordance with the UV–vis-NIR DRS (Fig. S6a). In addition, the time-resolved transient PL decay spectra recorded at the corresponding steady-state emission peaks reveal the shortened average lifetime ($\tau_{avg}$) of photogenerated charge carriers in Zn-PCN with respect to BCN (Fig. 4d and Table S3), meaning that the Zn single atoms...
atoms coordinated on PCN would effectively accelerate the charge transfer process [28].

Another solid evidence strongly supporting the improved charge transfer ability triggered by the introduced Zn single atoms could be accounted for by surface photovoltage (SPV) spectra. As shown in Fig. 4e, the SPV signals of both BCN and Zn-PCN are positive, a characteristic of n-type semiconductors, revealing that photogenerated holes transfer to the surface of PCN under illumination [44]. The more positive SPV signal observed for Zn-PCN than BCN indicates that...
the introduction of Zn single atoms would enable more holes to transfer to the surface of PCN [44]. This Zn single atoms promoted charge-carrier transfer process could be affirmed by the corresponding phase spectra (Fig. 4f), in which the phase changes from the first quadrant for BCN to the fourth quadrant for Zn-PCN [45]. Considering the electron-deficient feature of Zn single-atom sites, more photogenerated holes transfer to Zn single atoms under light irradiation, thus accelerating the charge transfer in Zn-PCN. Moreover, a new SPV signal appears at the wavelength range of 420–450 nm for Zn-PCN (Fig. 4e), again confirming the enhanced electron excitation originated from Zn-related midgap energy level (Fig. 4a) [16].

The charge transfer behavior in Zn-PCN, which was electrodeposited on a conductive substrate as a photoelectrode, was further evaluated by photoelectrochemical measurements. As revealed by the electrochemical impedance spectra of the as-deposited photoelectrodes (Fig. 4g), with Nyquist plots fitted by a typical equivalent circuit model (Table S4 and the inset in Fig. 4g), Zn-PCN exhibits much decreased charge transfer resistances in comparison with BCN, indicating the faster charge transfer processes in Zn-PCN [46], which is in accordance with the above PL and SPV results. These consistent conclusions can also be supported by the photoelectric response tests (Fig. 4h), which were recorded for several on–off cycles of intermittent irradiation. It is clear that the transient photocurrent densities of Zn-PCN are much higher than that of BCN, again providing reasonable evidence for the promoted charge transfer in Zn-PCN for surface water redox reaction under light irradiation [47].

### 3.5 Photocatalytic Performance for Water Splitting

The photocatalytic activities of the as-prepared samples were evaluated under visible-light irradiation ($\lambda > 420$ nm) for $\text{H}_2$ production from aqueous solution containing hole scavenger or pure water without any sacrificial agent. With triethanolamine (TEOA) used as hole scavenger, BCN could produce $\text{H}_2$ at a low rate of only $82.7 \mu\text{mol h}^{-1} \text{g}^{-1}$ (Fig. 5a, b). By comparison, the Zn-PCN samples exhibit gradually increased photocatalytic activity for $\text{H}_2$ production, depending on the increasing amounts of Zn single atoms. The highest $\text{H}_2$ production rate could reach $1172.9 \mu\text{mol h}^{-1} \text{g}^{-1}$ for Zn-PCN (4.79%), which is more than 14 times that of BCN (82.7 μmol h$^{-1}$ g$^{-1}$). It is generally accepted that BCN can hardly split pure water under visible light (Fig. 5c). Encouragingly, Zn-PCN(4.79%) is capable of producing $\text{H}_2$ and $\text{O}_2$ from pure water at a stoichiometric ratio of 2:1, with $\text{H}_2$ and $\text{O}_2$ production rates reaching 35.2 and 17.3 μmol h$^{-1}$ g$^{-1}$, respectively. To illustrate the origin of the $\text{O}_2$ product, $^{18}\text{O}$ isotope-labeled photocatalytic overall water ($\text{H}_2^{18}\text{O}$, 99%) splitting measurement was performed over Zn-PCN(4.79%). The molar content of the labeled $^{18}\text{O}_2$ with $m/z$ of 36 is detected to be 97.44% (Fig. 5d), which is very close to the theoretical value of 98.01% as calculated by the random adsorption model [16]. It could be then confirmed that $\text{O}_2$ is photocatalytically generated via pure water splitting. The photocatalytic stability was investigated with Zn-PCN(4.79%). As shown in Fig. S8, no noticeable degradation in photocatalytic activity was observed during the reaction for 72 h, implying that Zn-PCN is very stable for water splitting reactions. The almost unchanged XANES, XPS and FT-EXAFS spectra before and after the long-term photocatalytic reaction (Fig. S9) further demonstrate the good stability of Zn-PCN for overall water splitting.

### 3.6 Mechanism Analysis for Photocatalytic Overall Water Splitting

All the above experimental observations and theoretical calculations preliminarily infer that the enhanced photocatalytic performance achieved for Zn-PCN should be attributed to the Zn single-atom modulated electronic configuration enabling promoted electron excitation and accelerated charge separation/transfer, as well as the abundant surface active sites derived from the 3D porous structure for accelerated surface redox reactions. To deepen the insights into the underlying fundamentals determinative to the improved photocatalytic activity, water reduction and oxidation reaction paths were theoretically calculated for BCN and Zn-PCN, with free energy diagrams shown in Fig. 6a, b. For BCN, an energy difference of 0.82 eV is required for the formation of $\text{H}^+$ at C atoms for $\text{H}_2$ evolution (Fig. 6a). By comparison, Zn-PCN holds a much smaller energy difference of 0.58 eV (Fig. 6a), meaning the lower energy barrier for $\text{H}^+$ formation [41, 48]. As illustrated in the differential charge density calculations (Fig. 3h), rather than N atoms in BCN, the introduced Zn single atoms with an electron-deficient state are favorable to act as water oxidation active sites in
Zn‑PCN. For O₂ evolution, Zn‑PCN also has lower energy barriers than BCN from H₂O adsorption to intermediates (i.e., OH*, O* and OOH*) formation (Fig. 6b). These DFT calculation results indicate that water splitting into H₂ and O₂ at Zn‑PCN requires lower energy than at BCN, with more electrons and holes accumulated at the water reduction and oxidation active sites, respectively, leading to the enhanced adsorption and activation of H₂O molecules and then the reduced energy barriers for the formation of H* and OH* intermediates.

On the basis of the comprehensive analysis above in electronic structures, band structures, photoexcitation and charge transfer dynamics as well as water redox reaction paths, a possible mechanism for the improved photocatalytic activity over Zn‑PCN is shown in Fig. 6c. Along with the 3D porous structure beneficial to abundant surface active sites for water splitting, the anchored Zn single atoms featured with an electron‑deficient Zn‑N₆ configuration effectively modulate the electronic configuration of Zn‑PCN, which can induce a midgap energy level, increase the carrier concentration, strengthen the charge polarization in Zn‑PCN and redistribute the charges with electrons and holes enriched at the water redox active sites, contributing to promoted electron excitation, accelerated charge separation and transfer as well as reduced overpotentials of water redox reactions, finally enabling Zn‑PCN to realize visible‑light photocatalytic pure water splitting with H₂ and O₂ stoichiometrically evolved.

4 Conclusion

In this work, atomically dispersed Zn‑anchored 3D sponge‑like PCN (Zn‑PCN) characteristic of a unique Zn‑N₆ electron‑deficient configuration was successfully
synthesized via an intermediate coordination strategy. The obtained Zn-PCN achieves significantly improved photocatalytic H₂ production activity than BCN from water splitting in the presence of sacrificial agent, as well as realizes photocatalytic overall water splitting to stoichiometrically produce H₂ and O₂ with good durability under visible light. It is experimentally and theoretically revealed that such excellent photocatalytic water splitting performance should be benefitted from the modulated electronic configuration of PCN by the anchored Zn single atoms, which can induce a midgap energy level, increase the carrier concentration and strengthen the charge polarization, resulting in enhanced electron excitation, accelerated charge separation/transfer and reduced water redox barriers. In addition, the 3D porous structure would provide abundant surface reactive sites for water splitting. This work reports a novel single-atom photocatalyst for water splitting, and deepens the understanding of structure–activity relationship from the viewpoint of atomic coordination and electronic configuration, which is believed to provide informative principles for the design of efficient photocatalysts toward solar energy conversion.

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