The nature and descendants of Lyman-break galaxies in the Λ cold dark matter cosmology

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ABSTRACT

We predict the formation histories, properties and descendants of Lyman-break galaxies (LBGs) in the Λ cold dark matter cosmology. In our model, which incorporates a top-heavy initial mass function in starbursts, we find that most LBGs are starbursts triggered by minor mergers of galaxies. The duration of the LBG phase is predicted to be quite short, ∼20–60 Myr. We investigate the distributions of stellar and halo masses and morphologies for bright (L_{UV} > L_{UV}^*), and faint (L_{UV} > 0.1L_{UV}^*) LBGs at z = 3, 6 and 10 [where we classify LBGs according to their rest-frame ultraviolet (UV) luminosities relative the observed characteristic luminosity L_{UV}^* at z ≈ 3]. Bright LBGs at z = 3 are predicted to have median stellar masses ∼1 \times 10^9 h^{-1} M_☉ and host halo masses ∼3 \times 10^{11} h^{-1} M_☉, and be typically mildly disc dominated in stellar mass. On the other hand, faint LBGs at z = 10 are predicted to have median stellar masses of only ∼1 \times 10^7 h^{-1} M_☉ and host halo masses ∼2 \times 10^{10} h^{-1} M_☉, and be generally bulge dominated. Bright LBGs at z = 3 evolve to present-day galaxies with median stellar mass ∼5 \times 10^{11} h^{-1} M_☉ (comparable to the Milky Way), consisting of roughly equal numbers of disc- and bulge-dominated systems, and hosted by haloes with median mass ∼2 \times 10^{13} h^{-1} M_☉ (corresponding to medium-size galaxy groups). The model predicts that 40 per cent of Milky Way mass galaxies at the present day have a bright LBG progenitor in the redshift range 3 < z < 4, while 95 per cent have a faint LBG progenitor in the same redshift range and 7 per cent have a faint LBG progenitor at 10 < z < 11. With our multiwavelength model, we also investigate the overlap between the LBG population and that of selected sub-millimetre galaxies (SMGs); at z = 3, only ∼1 per cent of bright LBGs are also predicted to be bright SMGs (with an 850 μm flux in excess of 5 mJy).

Key words: galaxies: evolution – galaxies: formation – galaxies: high-redshift – galaxies: starburst – cosmology: theory.

1 INTRODUCTION

High-redshift galaxy populations identified using the Lyman-break drop-out technique have played a pivotal role in the development of our understanding of the global star formation history of the Universe (Madau et al. 1996; Steidel et al. 1999). The star formation rate (SFR) density inferred from measurements of the galaxy luminosity function in the rest-frame ultraviolet (UV) displays a broad peak around redshift ∼ 2, prior to a steep decline towards the present day (Lilly et al. 1996). This type of behaviour, which corresponds to a gradual build-up of the global stellar mass in galaxies, arises naturally in hierarchical structure formation models, as shown by Cole et al. (1994) and confirmed in a range of calculations with varying galaxy formation physics (e.g. Baugh et al. 1998; Somerville, Primack & Faber 2001; Springel & Hernquist 2003). In general terms, the SFR density is set by the abundance of haloes as a function of mass which are in place at a given epoch, the ability of gas to cool within these haloes and the effectiveness of supernovae at reheating gas which has cooled. The star formation density peaks around z ∼ 2 because gas can cool efficiently at high densities and stays in the cold phase because there is, for the first time, a large population of haloes in which supernova heating is ineffective. In this paper we investigate the properties of Lyman-break galaxies in a hierarchical galaxy formation model and show how they are related to galaxies in the present-day Universe.

The Lyman-break technique is based on observing a break bluewards of the Lyman limit at 912 Å in the rest-frame spectral energy

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distribution (SED) of a galaxy. The flux at these wavelengths is expected to be small for a number of reasons: (i) few photons can be produced by stars at these short wavelengths, (ii) such energetic photons can be absorbed by the atomic hydrogen in the interstellar medium within the galaxy and (iii) the photons can also be absorbed by intergalactic neutral hydrogen along the line of sight from the galaxy to the Earth. Galaxies identified with this technique are known as Lyman-break galaxies (LBGs; Steidel & Hamilton 1992; Steidel et al. 1996). The Lyman-break technique uses multiband photometry. By using a suitable set of filters, galaxies in a particular range of redshifts can be isolated. Ground-based observations are restricted to applying the Lyman-break technique for galaxies at $z > 2.5$. For this redshift range, the absorption edge, which occurs in the far-UV, is shifted to wavelengths that are not blocked by the Earth’s atmosphere. The multiband photometry is also used to make additional colour selections to minimize the contamination from foreground objects (both stars and lower redshift galaxies). LBGs are selected to be red in the shorter wavelength colour, due to the presence of the break, and blue in the longer wavelength colour, to show that they are star-forming galaxies. Hot, high-mass stars with short lifetimes tend to dominate the emission in the UV. Such stars spend a significant fraction of their lifetime within dense giant molecular clouds and their light is heavily extinct. Hence in order to infer a galaxy’s SFR from its UV luminosity an uncertain and quite often large correction for dust extinction is required. Furthermore, an assumption about the form of the initial mass function (IMF) with which stars are produced is also required in order to infer the SFR.

The nature of the Lyman-break population in hierarchical cosmologies has been probed using both semi-analytical modelling and gas dynamic simulations (Baugh et al. 1998; Somerville et al. 2001; Nagamine et al. 2005, 2010; Finlator et al. 2006; Night et al. 2006; Lo Faro et al. 2009). The semi-analytical approach was first applied to LBGs by Baugh et al. (1998), who showed that the model of Cole et al. (1994) could reproduce the observed surface density of LBGs (e.g. Steidel et al. 1996). Somerville et al. (2001) applied similar techniques to argue that bursts triggered by minor and major galaxy mergers played a role in producing the LBG population. The importance of galaxy mergers in LBGs is however still open to debate. Guo & White (2009) studied the properties and clustering of LBGs in the Munich semi-analytical galaxy formation model, and argued that LBGs are powered by quiescent star formation in galactic discs rather than by starbursts. These authors exploited the semi-analytical model to connect the galaxies identified at high redshift to their present-day descendants to study the properties of the descendants of LBGs.

In this paper we use the version of the Durham GALFORM galaxy formation simulation code published by Baugh et al. (2005) to study LBGs and their descendants (see Baugh 2006 and Benson 2010 for reviews of semi-analytical modelling). This model invokes a different form for the stellar IMF in merger-driven starbursts, in order to reproduce the observed counts and redshift distribution of galaxies measured at submillimetre wavelengths (submillimetre galaxies, SMGs). Baugh et al. (2005) showed that this model also matches the observed LBG luminosity function at $z = 3$, as well as giving a reasonable match to the present-day galaxy population. The adoption of a different IMF in merger-driven starbursts is controversial. Some authors have argued that it is not necessary to invoke a change to the form of the IMF to match the observed counts of SMGs (Hayward et al. 2011). We fully agree with this conclusion. However, the challenge is to not only match the counts but also reproduce the redshift distribution of SMGs and the properties of the local galaxy population. This broader objective has not been demonstrated in a model with a universal IMF and a self-consistent treatment of the effect of dust on the SED of galaxies.

The relaxation of the hypothesis of a universal IMF is contentious and so it is necessary to test as many other predictions of the model as possible. A clear strength of the semi-analytical approach is that multiwavelength predictions can be generated readily at any epoch. Furthermore, populations can be connected at different epochs, allowing the properties of the descendants of high-redshift galaxies to be studied. González et al. (2009) extended the local tests of this model to include the size–luminosity relation for disc-dominated galaxies. In González et al. (2011) we followed the evolution of SMGs and their descendants. In this paper, we carry out a similar analysis in the UV by studying the evolution and descendants of LBGs. This is a companion paper to Lacey et al. (2011), who presented the model predictions for the evolution of the far-UV luminosity functions over the redshift range $z = 3–20$, making a detailed comparison of these predictions with current observational data. Lacey et al. (2011) also made predictions for the amount of dust extinction and for the evolution of a wide variety of other properties of LBGs over a range of redshifts. The present paper instead looks in more detail at the formation histories and lifetimes of LBGs and the distributions of stellar and halo masses and morphologies for LBGs at selected redshifts and luminosities, and investigates the corresponding properties for their present-day descendants.

This paper is laid out as follows. In Section 2, we outline the GALFORM galaxy formation model. In Section 3, we review the model predictions for the luminosity function of LBGs, show examples of galaxy merger histories which include LBGs and investigate the triggering and duration of the LBG phase. We also present the probability that a present-day galaxy or halo had an LBG progenitor. In Section 4, we study different properties of the LBGs, such as stellar and host halo masses, and morphologies at redshifts $z = 3$, 6 and 10, and discuss the dominant mode of star formation behind the LBGs. In Section 5, we examine the overlap between the LBG and SMG populations. Finally, we present a summary in Section 6.

We quote magnitudes on the AB system and define the Hubble parameter as $h = H_0/100$ km s$^{-1}$ Mpc$^{-1}$.

## 2 THE GALAXY FORMATION MODEL

We first give a brief overview of the Durham semi-analytical galaxy formation code, GALFORM, before discussing two aspects which are of particular relevance for this paper, galaxy mergers and the form of the IMF adopted in different modes of star formation. Full details of GALFORM can be found in Cole et al. (2000) and Benson et al. (2003). The model used here is described in full in Baugh et al. (2005) and Lacey et al. (2008).

### 2.1 Basic components

Galaxies are assumed to form inside dark matter haloes, with their subsequent evolution being controlled by the merger and accretion history of the halo (White & Rees 1978). The physical processes modelled in GALFORM include: (i) the hierarchical assembly of dark matter haloes; (ii) the shock heating and virialization of gas inside the gravitational potential wells of dark matter haloes; (iii) the radiative cooling of the gas to form a galactic disc; (iv) star formation in the cool gas; (v) the heating and expulsion of cold gas through feedback processes such as stellar winds and supernovae; (vi) chemical evolution of gas and stars; (vii) galaxy mergers within a common dark halo as a result of the decay of a galaxy’s orbit.
due to dynamical friction; (viii) the evolution of the stellar populations using population synthesis models and (ix) the extinction and reprocessing of starlight by dust.

We use the same galaxy formation parameters as adopted by Baugh et al. (2005) and Lacey et al. (2008). The model is set in a cold dark matter plus cosmological constant cosmology, with the present-day matter density parameter \( \Omega_0 = 0.3 \), a cosmological constant \( \Omega_\Lambda = 0.7 \), a baryon density \( \Omega_b = 0.04 \) and a power spectrum normalization given by \( \sigma_8 = 0.93 \).

In this paper, one of our primary goals is to connect high-redshift galaxies with their present-day descendants. This requires a prescription for building halo merger histories which is accurate over a long interval in time. To achieve this we use the Monte Carlo algorithm of Parkinson, Cole & Helly (2008), which is a modified version of that proposed by Cole et al. (2000) and used in Baugh et al. (2005). The Parkinson et al. recipe has been tuned to match the merger histories extracted from the Millennium Simulation (Springel et al. 2005; see Harker et al. 2006 for a description of the construction of merger trees from the simulation). We note that Baugh et al. did not make any predictions which depended on the accuracy of the trees over long time intervals. Rather, they laid down grids of haloes at a range of redshifts to compute the luminosity function of SMGs and other galaxies at these redshifts, and then, by integration, the number counts of SMGs and the cosmic star formation history. We have checked that using the Parkinson et al. tree algorithm does not alter the predictions presented by Baugh et al.

Dust extinction is expected to have a large effect on the far-UV luminosities of LBGs. We calculate dust extinction in the model using the approach described in Lacey et al. (2011), and we refer the reader to that paper for a discussion of the effects of dust extinction on LBGs. In summary, the dust extinction is calculated using a two-component dust model, with molecular clouds within which stars form embedded in a diffuse dust medium mixed with the stars and distributed in an exponential disc. All of the predictions which we show in this paper include the effects of dust extinction. The effect of dust extinction on the slope of the UV continuum in the model is discussed in Gonzalez-Perez et al. (in preparation).

### 2.2 Galaxy mergers and merger-driven starbursts

When dark matter haloes merge, the galaxy in the most massive halo is assumed to become the central galaxy in the new halo while the other galaxies present become its satellites. In the standard form of GALFORM, we assume that central and satellite galaxies can form stars from their cold gas reservoirs, but that only central galaxies can continue to accrete cold gas. (Font et al. 2008 explore a variant of this model in which satellites are allowed to retain some fraction of their hot halo, depending on their orbit within the main halo, with gas allowed to cool on to the satellite discs.)

In the standard version of the model, mergers are assumes to happen only between satellites and central galaxies within the same halo (although Kim et al. 2009 considered a variant of this). Bursts of star formation are assumed to be triggered by galaxy mergers. Two types of mergers are defined, major mergers and minor mergers, according to whether the ratio of the mass of the smaller to that of the larger galaxy exceeds the value of the parameter \( f_{\text{flip}} \), i.e. \( M_2/M_1 > f_{\text{flip}} \). Bursts are assumed to be triggered in all major mergers, and also in minor mergers which satisfy \( M_2/M_1 > f_{\text{burst}} \) and where the ratio of cold gas mass to stellar plus cold gas mass in the larger galaxy exceeds \( f_{\text{gas, crit}} \). The parameter values adopted in the Baugh et al. model are \( f_{\text{flip}} = 0.3, f_{\text{burst}} = 0.05 \) and \( f_{\text{gas, crit}} = 0.75 \). In a major merger, the stellar discs of the merging galaxies are assumed to be transformed into a new stellar spheroid (resulting in a pure bulge galaxy) and any gas present in the discs is converted into spheroid stars in a burst. In minor mergers, the stellar disc of the larger galaxy is preserved and the stars of the smaller, accreted galaxy are added to its bulge (resulting in a disc plus bulge galaxy). If the condition for a burst in a minor merger is satisfied, then the gas of the smaller galaxy is converted into stars. If the condition is not satisfied, any accreted gas is added to the gas disc of the central galaxy. From this description of the possible scenarios following a galaxy merger, it is readily apparent that mergers change the morphology of galaxies in the model, driving the morphology towards early types (bulge dominated), with the cooling of gas moving the galaxy towards late types (disc dominated).

### 2.3 The stellar IMF in quiescent and burst modes of star formation

An important feature of the Baugh et al. model is that stars are assumed to form with different stellar IMFs in different modes of star formation. Stars either form quiescently in galactic discs or in bursts triggered by some galaxy mergers, as explained in the subsection above. In the case of quiescent star formation, a standard solar neighbourhood IMF is adopted, following the form proposed by Kennicutt (1983), where the number of stars produced depends on mass as \( \text{d}N/\text{d}m \propto m^{-x} \), with \( x = 0.4 \) for \( m < 1 \, M_\odot \) and \( x = 1.5 \) for \( m > 1 \, M_\odot \). In the case of merger-driven starbursts, a top-heavy IMF is assumed where \( x = 0 \). The use of a top-heavy rather than solar neighbourhood IMF has three consequences: (i) it increases the amount of radiation emitted in the UV, as a larger proportion of massive stars are produced, (ii) there is higher yield of metals \( (p = 0.15 \) compared with \( p = 0.023 \) for the Kennicutt IMF), leading to more dust, and (iii) the fraction of mass turned into stars which is recycled to the interstellar medium is larger (\( R = 0.91 \) compared with \( R = 0.41 \) for the Kennicutt IMF). For a discussion of observational support for a top-heavy IMF see Lacey et al. (2008, 2010).

Finally, in closing this section we remind the reader that an alternative galaxy formation scenario was presented by Bower et al. (2006), who invoke a different mechanism to suppress the formation of galaxies in massive haloes than that used by Baugh et al. Bower et al. also invoke a standard IMF in both modes of star formation. Whilst this model gives a good match to the observed evolution of the \( K \)-band luminosity function and the inferred stellar mass function, it greatly underpredicts the number of SMGs and gives a poorer match to the LBG luminosity function than the Baugh et al. model (see Lacey et al. 2011). For further discussion of the differences between these two models, see Almeida, Baugh & Lacey (2007), Gonzalez-Perez et al. (2009) and González-Perez et al. (2009).

### 3 LBG FORMATION HISTORIES AND THE PROBABILITY OF HAVING A LBG PROGENITOR

In this section, we present some basic predictions for LBGs in the Baugh et al. model. We first compare the model predictions for the luminosity function in the UV at \( z = 3, 6 \) and 10 with observations, and use this to set up definitions of bright and faint LBGs. (For a more detailed comparison with observed LBG luminosity functions see Lacey et al. 2011.) We then present some illustrative galaxy formation histories for present-day galaxies which had LBG progenitors. Finally, we show the probability that a present-day

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galaxy had an LBG progenitor and that a present-day halo hosted an LBG.

### 3.1 The UV luminosity function: defining bright and faint Lyman-break galaxy samples

As outlined in Section 1, the Lyman-break technique relies on colour–colour selection to identify galaxies within a particular redshift range, which is set by the redshifted Lyman-break spectral feature falling between two of the filters used to image the galaxies. In the model we can make predictions for the galaxy population at any desired redshift by construction and so there is no need to apply any colour selection. Hence, we will use the far-UV (1500 Å) luminosity as our criterion to identify samples of LBGs. We focus on three redshifts, \( z = 3, 6 \) and 10, which are representative of the range covered by current LBG studies (Steidel et al. 2003; Bouwens et al. 2007, 2011).

Baugh et al. (2005) presented predictions for the rest-frame UV luminosity function at \( z = 3 \). Here we revisit this comparison, adding new observational data at \( z = 3 \) and going to \( z = 10 \) (see also Lacey et al. 2011). In Fig. 1 we show the predicted rest-frame 1500 Å luminosity function at \( z = 3, 6 \) and 10, including the effects of dust extinction, calculated in a self-consistent way using the predicted scalelengths of the disc and bulge components and the metallicity of the cold gas. The model matches the observational data at \( z = 3 \), to within the scatter between data sets, is in good agreement with the measurement by Bouwens et al. (2007) at \( z = 6 \) and is consistent with the tentative measurement and upper limits estimated by Bouwens et al. (2011) at \( z = 10 \). Following the common practice in observational studies, we define a characteristic luminosity \( L_{\mathrm{UV}}^{*} \) using the observed position of the break in the UV luminosity function at \( z = 3 \). We define a galaxy as a bright LBG if its UV luminosity is brighter than \( L_{\mathrm{UV}}^{*} \) and as a faint LBG if its UV luminosity exceeds \( 0.1L_{\mathrm{UV}}^{*} \) (a sample which includes the bright LBG sample, but which is dominated in number by galaxies close to the luminosity cut). We apply this definition at all redshifts. Note that \( L_{\mathrm{UV}}^{*} \) is taken to be the characteristic luminosity at \( z = 3 \) found by Steidel et al. (1999) (corresponding to \( M_{\mathrm{UV}}^{*} = -20.3 + 5 \log h \) in our cosmology). Fig. 1 shows the location of the 0.1\( L_{\mathrm{UV}}^{*} \) and \( L_{\mathrm{UV}}^{*} \) limits using vertical lines.

We note in passing that there is substantial evolution in both the observed and predicted UV luminosity functions between \( z = 3 \) and 6. The observed characteristic luminosity \( L_{\mathrm{UV}}^{*}(z=6) \) at \( z = 6 \) (\( M_{\mathrm{UV},z=6}^{*} = -19.5 + 5 \log h \) in our cosmology; Bouwens et al. 2007) is nearly a magnitude fainter than the \( z = 3 \) value of \( L_{\mathrm{UV}}^{*} \). The number density of galaxies seen at the \( z = 3 \) value of \( L_{\mathrm{UV}}^{*} \) drops by around a factor of 5 between \( z = 3 \) and 6. At \( z = 10 \), the abundance of galaxies with this luminosity is predicted to be several hundred times lower than at \( z = 3 \). A comprehensive study of the evolution of the LBG luminosity function, including a more detailed comparison with observational data, is presented in Lacey et al. (2011).

![Figure 1](http://mnras.oxfordjournals.org/)

**Figure 1.** The galaxy luminosity function in the rest-frame UV (1500 Å) at \( z = 3 \) (top panel), \( z = 6 \) (middle panel) and \( z = 10 \) (bottom panel). The symbols with error bars show observational estimates taken from Arnouts et al. (2005) and Reddy & Steidel (2009) at \( z = 3 \), Bouwens et al. (2007) at \( z = 6 \) and Bouwens et al. (2011) at \( z = 10 \) as indicated by the legends. The predicted luminosity functions include the effect of dust extinction and the observational estimates have not been corrected to remove extinction. The vertical red lines indicate \( L_{\mathrm{UV}}^{*} \) and 0.1\( L_{\mathrm{UV}}^{*} \) (defined at \( z = 3 \)) which we use to construct samples of bright and faint LBGs, respectively.

### 3.2 Example of formation histories of model galaxies with Lyman-break progenitors

As examples of the different galaxy formation and merger histories which can produce bright LBGs, we plot in the left-hand panels of Fig. 2 galaxy merger trees for three present-day galaxies. The mass of the galaxies increases down the page. The trees are constructed by running GALFORM with many output redshifts. We stop plotting the trees at \( z = 6 \) or when a branch falls below a minimum galaxy mass (in stars and cold gas) of \( 10^{6} h^{-1} \) M\(_{\odot} \). A branch or trunk of the tree is plotted at each output redshift using a circle. The size of the circle is proportional to the stellar mass of the galaxy and the colour reflects the type of galaxy: green for a faint LBG, red for a bright LBG and blue otherwise.

The galaxy trees are plotted in the following way. Stepping back in time from \( z = 0 \), at each merger we plot the most massive branch on the left and the other branches to the right of this. The main progenitor branch is hence the leftmost plotted branch traced back from \( z = 0 \) in this way. Note that at \( z > 0 \), the main progenitor branch does not necessarily represent the most massive progenitor across the whole tree at a given epoch. The largest progenitor at any time could jump from one branch of the galaxy tree to another, so its selection would not necessarily define a smooth path back in redshift. We have chosen to plot examples which have at least one LBG in the most massive (leftmost) progenitor branch.

In the first example (Fig. 2, top-left), we show the galaxy merger tree for a galaxy which at \( z = 0 \) has a stellar mass of \( M_{*} = 2.5 \times 10^{11} h^{-1} \) M\(_{\odot} \) and has a bright LBG progenitor at \( z = 6 \). In the second example (Fig. 2, middle-left), we show a galaxy with \( M_{*} = 6.6 \times 10^{11} h^{-1} \) M\(_{\odot} \) at \( z = 0 \) with a bright LBG progenitor at both \( z = 6 \) and 10.
Figure 2. Left-hand panels: galaxy merger trees for different $z = 0$ galaxies that have at least one LBG in the most massive progenitor branch. From top to bottom, the stellar masses of the galaxies at $z = 0$ are $M_* = 2.5 \times 10^9, 6.6 \times 10^{10}$ and $2.1 \times 10^{11} h^{-1} M_\odot$, respectively. The size of the symbol is proportional to the stellar mass as indicated by the key in the top-left panel. The green circles are drawn for galaxies with $L_{UV} > L_*^{UV}$, the red circles for $L_{UV} > L_*^{UV}$ and the blue if $L_{UV} < L_*^{UV}$. Right-hand panels: evolution with redshift of the galaxy on the main progenitor branch. Top subpanel: evolution of the bulge-to-total stellar mass ratio ($B/T$). The dashed vertical line indicates the redshift at which the central galaxy (black line) becomes a satellite galaxy (plotted using a green line) after a halo merger. Main right-hand subpanel: the symbols show the evolution of the stellar mass (blue), cold gas mass (green), host dark matter halo mass (black) and the stellar mass formed in ongoing bursts (red) in units according to the left-hand label. The red line shows the evolution of the ratio $L_{UV}/L_*^{UV}$ with redshift, in units according to the right-hand axis. Note that $L_*^{UV}$ is the observationally determined reference value at $z = 3$ (Steidel et al. 1999).
and 3. In the third example (Fig. 2, bottom-left), we show the galaxy merger tree for a galaxy with $M_* = 2.1 \times 10^{11} \, h^{-1} \, M_\odot$ at $z = 0$ with a bright LBG progenitor at $z = 6$. Note that these trees are purely illustrative examples, chosen to show the range of complexity of the trees, and are not intended to be a statistically representative sample.

The right-hand panels of Fig. 2 show the evolution of various properties of the galaxy in the most massive progenitor branch. The upper inset or subpanel in each case shows $B/T$. If present, a vertical line in this panel marks the epoch when the most massive progenitor ceased to be a central galaxy and became a satellite. In the examples plotted here, this only happens in the case of the least massive galaxy tree shown (upper-right, Fig. 2), at $z \approx 1.1$. The main right-hand panel shows the UV luminosity of the most massive progenitor, in units of $L_{UV}$ (red line; see the right axis for units) and its stellar mass (blue symbols), cold gas mass (green symbols), stellar mass formed in bursts (red symbols) and host dark halo mass (black symbols; see the left axis for units). In the case of the least massive galaxy (Fig. 2, top-right), the most massive progenitor experiences a burst at $z = 6$ which makes it a bright LBG. The stellar mass added to the bulge during the burst increases $B/T$, and the cold gas supply is depleted. Around $z \approx 5.2$ there is another burst of star formation, triggered by a major merger, which leads to a large increase in the bulge-to-total ratio of the galaxy from 0.3 to 0.9. However, there is less star formation associated with this burst, due to the depletion of the gas reservoir prior to the burst, and the most massive progenitor is only a faint LBG. Another burst happens at $z \approx 3.5$, which makes the most massive progenitor galaxy a faint LBG but with relatively little change in $B/T$. $B/T$ gradually declines as quiescent star formation adds mass to the disc. Finally, after a steady increase in mass, the host halo is accreted by a much more massive halo at $z = 1.1$, when the most massive progenitor becomes a satellite galaxy. The present-day galaxy is a satellite in a halo of mass $10^{14} \, h^{-1} \, M_\odot$. The next most massive galaxy (middle panels of Fig. 2) displays two epochs at which the most massive progenitor is a bright LBG ($z \approx 6$ and 3.6). The most massive progenitor branch in the most massive galaxy (bottom panels of Fig. 2) again shows several instances when the progenitor is a bright LBG. It undergoes a major merger at $z \approx 1.5$ which increases the stellar mass by 30 per cent (barely visible in the figure), triggers a bright LBG phase and changes the morphology from $B/T \approx 0$ to 1.

It is interesting to note that in all cases shown, the stellar mass of the most massive progenitor branch (shown by the blue points in the right-hand panels) increases steadily with redshift, without displaying any sudden large jumps. This implies that the bursts of star formation in these examples are not responsible for large changes in the stellar mass of the galaxy. This is a result in part of the large recycled fraction ($R = 0.91$) for stars forming with the top-heavy IMF, together with the strong feedback from supernova explosions.

### 3.3 The triggering and duration of the Lyman-break phase

We can use the formation histories of model LBGs to answer some important questions. The first of these is: what is the mechanism of the star formation responsible for the galaxy being seen as an LBG at a particular redshift? In the present model, the possibilities are: a starburst triggered by a minor galaxy merger; a starburst triggered by a major galaxy merger and quiescent star formation in a galaxy disc. To determine which of these dominates in a particular case, we move back in time through the galaxy formation history, starting at the redshift at which a galaxy is identified as an LBG.

Figure 3. The triggering mechanism of the LBG phase. For each model LBG, we trace its history back in time from the redshift at which it is identified, to determine whether the UV luminosity is dominated by stars formed in a burst triggered by a minor merger (blue), a burst triggered by a major merger (red) or stars formed quiescently (green). The top panel shows the fractions for these different triggering mechanisms for bright LBGs ($L_{UV} > L_{UV}^* \, 0.1$) identified at $z = 3$ and 6, while the bottom panel shows the fractions for faint LBGs ($L_{UV} \, 0.1L_{UV}^*$) identified at $z = 3, 6$ and 10. (No bright LBGs are predicted at $z = 10$.)

If at any time during the current LBG phase the burst and quiescent SFRs satisfy the condition $SFR_{quiescent} > 3 \, SFR_{burst}$, then we classify the LBG phase as being produced by quiescent star formation; otherwise we find the first galaxy merger (going back in time) and identify whether it is a minor or major merger. We include the factor 4 in the condition on the SFRs because stars forming with the top-heavy burst IMF produce roughly four times higher UV luminosities than the same mass of stars formed with the quiescent IMF. The results of this classification exercise are shown in Fig. 3, for bright LBGs ($L_{UV} > L_{UV}^*$) identified at $z = 3$ and 6 (top panel), and for faint LBGs ($L_{UV} \, 0.1L_{UV}^*$) identified at $z = 3, 6$ and 10 (bottom panel). The fractions shown there are calculated for a volume-limited sample of LBGs at the identification redshift. We see from the figure that LBGs are predicted to be dominated by starbursts triggered by minor galaxy mergers (around 60–80 per cent of cases) in all of the cases examined here. Minor mergers dominate major mergers by a factor of 2–4. Quiescent star formation is predicted to account for at most a few per cent of LBGs, with the exception of faint LBGs at $z = 3$, where this fraction increases to around 20 per cent. We also see that the quiescent fraction decreases with increasing redshift for both bright and faint LBGs, and is much smaller for bright LBGs than for faint LBGs at the same redshift. These results about the dominance of bursts over quiescent star formation in LBGs are in agreement with the analysis by Lacey et al. (2011).
a statistical study and compute the probability that a present-day galaxy had a progenitor that was an LBG or that a present-day halo had a progenitor halo which hosted an LBG.

The predictions for the fraction of present-day galaxies that are descendants of LBGs are presented in Fig. 5, with the top panels showing results for descendants of bright LBGs at $3 < z < 4$ and $6 < z < 7$, and the lower panels results for descendants of faint LBGs at $3 < z < 4$, $6 < z < 7$ and $10 < z < 11$. (The model does not predict a significant number of bright LBGs at $z = 10$.) In the left-hand panels we plot the fraction of present-day galaxies, as a function of their stellar mass, which are predicted to have at least one bright LBG progenitor in the specified redshift range. In the right-hand panels we show the fraction of present-day haloes having progenitor haloes which hosted an LBG in the specified redshift range. We see that the probability of having an LBG progenitor increases with increasing present-day mass (both stellar and halo), but decreases with increasing LBG redshift.

Considering first the bright LBGs, we find that galaxies with stellar masses of $6.3 \times 10^{10}$, $6.1 \times 10^{10}$ and $1.8 \times 10^{11} h^{-1} M_\odot$ are predicted to, respectively, have 10, 50 and 90 per cent probabilities of having at least one bright LBG progenitor in the range $3 < z < 4$. The corresponding stellar masses for bright LBGs at $6 < z < 7$ are 4–15 times larger, being $9.0 \times 10^{10}$, $2.7 \times 10^{11}$ and $1.3 \times 10^{12} h^{-1} M_\odot$, respectively. A Milky Way mass galaxy (4 $\times$ $10^{10}$ $h^{-1} M_\odot$; Binney & Tremaine 2008) is predicted to have a 40 or 3 per cent probability of having a bright LBG progenitor at $3 < z < 4$ or $6 < z < 7$, respectively. The right-hand panels of Fig. 5 show that 50 per cent of present-day haloes with a mass greater than $4.3 \times 10^{12} h^{-1} M_\odot$ are predicted to have halo progenitors that hosted bright LBGs at $3 < z < 4$. The corresponding mass for bright LBGs at $6 < z < 7$ is an order of magnitude larger, $4.8 \times 10^{13} h^{-1} M_\odot$.

Considering now the faint LBGs, we find that 10, 50 and 90 per cent of galaxies with present-day stellar masses of $1.3 \times 10^{10}$, $2.6 \times 10^{9}$ and $2.2 \times 10^{10} h^{-1} M_\odot$ are predicted to have had faint LBG progenitors at $3 < z < 4$. These masses increase steeply to $1.1 \times 10^{10}$, $1.6 \times 10^{10}$ and $7.7 \times 10^{10} h^{-1} M_\odot$ for faint LBGs at $6 < z < 7$, and steeply again to $5.3 \times 10^{10}$, $2.0 \times 10^{11}$ and $4.2 \times 10^{11} h^{-1} M_\odot$ at $10 < z < 11$. A Milky Way mass galaxy is predicted to have a 95 or 70 per cent probability of having a faint LBG progenitor at $3 < z < 4$ or $6 < z < 7$, respectively. This probability drops to 7 per cent for a $10 < z < 11$ faint LBG progenitor. Expressed in terms of halo masses (right-hand panels of Fig. 5), we find that 50 per cent of haloes today with masses larger than $2.6 \times 10^{11} h^{-1} M_\odot$ are predicted to have halo progenitors that hosted faint LBGs at $3 < z < 4$, while for $6 < z < 7$ this increases to $1.4 \times 10^{12} h^{-1} M_\odot$. Only for masses greater than $2.7 \times 10^{13} h^{-1} M_\odot$ are 50 per cent of haloes predicted to have had at least one faint LBG progenitor at $10 < z < 11$.

4 THE PROPERTIES OF LBGs AND THEIR DESCENDANTS

Having defined samples of bright and faint LBGs, we now compare the model predictions for their stellar mass, host halo mass and morphologies at redshifts $z = 3$, 6 and 10. The semi-analytical model allows us to follow the subsequent evolution of the LBGs, so we also present the predictions for the analogous properties of the descendants of the LBGs. We also consider what is the dominant mode of star formation — burst or quiescent — in LBGs. Note that at $z = 10$, the model does not predict a significant population of bright galaxies.

### 3.4 The fraction of present-day galaxies and haloes with an LBG progenitor

Having shown some illustrative examples of galaxy merger trees with LBG progenitors in the previous subsection, we now perform
Figure 5. The predicted fraction of present-day galaxies (left) and haloes (right) with either bright LBG progenitors (i.e. $L_{UV} > L_{\ast, UV}$) in the top set of panels, or faint LBG progenitors (i.e. $L_{UV} > 0.1L_{\ast, UV}$) in the bottom set of panels. Left-hand panels: the fraction of present-day galaxies, as a function of their stellar mass, which are predicted to have one or more LBG progenitors. Right-hand panels: the fraction of present-day haloes, as a function of mass that are predicted to have hosted at least one LBG in a progenitor halo. Fractions are shown for LBG progenitors at $3 < z < 4$, $6 < z < 7$ and $10 < z < 11$ (the last for faint LBGs only), as labelled. Descendant galaxies are separated into bulge dominated ($B/T > 0.5$, red) and disc dominated ($B/T < 0.5$, blue). The black line corresponds to the case of all LBG descendants and is the sum of the blue and red lines.

4.1 Stellar masses

The model predicts that the stellar masses of LBGs identified at a given UV luminosity decrease with increasing redshift. On the other hand, LBGs (see Fig. 1). When we construct the distribution of properties of the descendants of LBGs, each descendant galaxy is counted once (with a weight depending on its abundance), regardless of the number of LBG progenitors it possesses.
hand, the stellar masses of their present-day descendants show the opposite trend, increasing with increasing redshift. This behaviour is seen for both bright and faint LBGs, as we now discuss in detail.

We plot the stellar mass distributions of LBGs and of their present-day descendants in the left- and right-hand panels of Fig. 6. The top panels show results for bright LBGs ($L_{UV} > L_{UV}^\text{10}$) at $z = 3$ and 6, while the lower panels show results for faint LBGs ($L_{UV} > 0.1L_{UV}^\text{10}$) at $z = 3$, 6 and 10.

We see that the median stellar mass of bright LBGs at $z = 3$ is $M_\star = 1.3 \times 10^{10} h^{-1} M_\odot$, which is six times larger than that of their faint counterparts at this redshift and $\approx 2.5$ times larger than the median mass of bright LBGs at $z = 6$. The median stellar masses of the z = 0 descendants (right-hand panels) of these two LBG populations ($z = 3$ and 6) are predicted to be $M_\star = 5 \times 10^{10} h^{-1} M_\odot$ for the bright LBGs at $z = 3$ and $M_\star = 1 \times 10^{11} h^{-1} M_\odot$ for the bright LBGs at $z = 6$. We separate the galaxies into bulge-dominated (if $B/T > 0.5$) and disc-dominated galaxies (i.e. $B/T < 0.5$). Bright LBGs at $z = 3$ and 6 are predicted to be typically (60 per cent) disc dominated. However, while the descendants of bright LBGs at $z = 3$ are also mainly disc dominated, descendants of bright LBGs at $z = 6$ are mainly bulge dominated.

The median stellar mass of faint LBGs at $z = 3$ is $M_\star = 2 \times 10^{10} h^{-1} M_\odot$, which is an order of magnitude larger than the median mass of faint LBGs at $z = 6$. Faint LBGs at $z = 10$ are 25 times less massive than faint LBGs at $z = 3$. The median stellar masses of the descendants of the $z = 3$ and 6 populations are predicted to be similar ($M_\star \approx (3-5) \times 10^{11} h^{-1} M_\odot$). The median descendant mass of a $z = 10$ faint LBG is an order of magnitude larger than for faint LBGs identified at $z = 3$ or 6. The descendants of LBGs show a broader distribution of stellar mass compared with the LBGs themselves. We also find that faint LBGs at $z = 6$ and 10 are predominantly bulge-dominated systems (75 per cent are predicted to have $B/T > 0.5$). At $z = 3$ there is a more balanced population, with almost half of the faint LBGs predicted to be disc-dominated systems. The descendants of faint LBGs identified at $z = 3$ and 6 have similar morphological mixes. They are predicted to be mainly disc-dominated systems (with 70 per cent having $B/T < 0.5$). The descendants of $z = 10$ faint LBGs are equally likely to be disc or bulge dominated. Among the descendants, those of lower mass ($M \lesssim 10^{11} h^{-1} M_\odot$) are typically predicted to be disc-dominated systems, whereas higher mass LBG descendants ($M \gtrsim 10^{11} h^{-1} M_\odot$) are mostly bulge-dominated systems in the model.

Stellar masses of LBGs have been estimated photometrically in a number of observational studies (e.g. Papovich, Dickinson & Ferguson 2001; Shapley et al. 2001; Verma et al. 2007; Stark et al. 2009). These are all based on fitting model SEDs to multicolour photometry. As discussed in Lacey et al. (2011), these photometric mass estimates are fraught with difficulty, since they depend sensitively on the assumed IMF, and also on assumptions about dust extinction, star formation history and metallicity. Since our model galaxies contain a mixture of stars formed with quiescent and burst IMFs, while the observational estimates all assume a solar neighbourhood (typically Salpeter) IMF, a direct comparison between stellar masses predicted by the model and values estimated observationally from photometric data is not meaningful. Instead, Lacey et al. made a direct photometric comparison between the models and observations. They plot the observer-frame flux at 3.6 \mu m against rest-frame far-UV flux, and find, using the same model as we do here, that the predictions come to within a factor of 3 of matching the observed relation found by Stark et al. (2009) for LBGs at $z = 4$–6. At the redshifts of the LBG samples, the 3.6 \mu m flux probes the optical rest frame, which is determined by the recent star formation history of a galaxy as well as its stellar mass. We will investigate this issue in more detail in a future paper.

As discussed in Lacey et al. (2011), our model predicts that LBGs should typically have large gas fractions. Unfortunately for testing the model, there are currently few observational estimates of gas masses of LBGs for the redshift range $z \approx 3$ which we are studying here, due to their molecular line emission typically being too faint for detection [although this situation should soon improve with Atacama Large Millimeter Array (ALMA) observations]. Two $z \approx 3$ LBGs having $L_{UV} \approx L_{UV}^\text{10}$ have been detected in CO emission, due to their being gravitationally lensed by factors of about 30 (Baker et al. 2004; Coppin et al. 2007; Riechers et al. 2010). For these two galaxies, Riechers et al. estimate total molecular gas masses from their CO emission of $6.6 \times 10^{9}$ and $3.3 \times 10^{9} h^{-1} M_\odot$, respectively, i.e. a mean of $5 \times 10^{9} h^{-1} M_\odot$. In contrast, our model predicts a median gas mass (including both molecular and atomic gas) of $\approx 5 \times 10^{9} h^{-1} M_\odot$ for the same redshift and UV luminosity, with roughly a factor 2 spread around this value. However, the Riechers et al. estimate assumes a conversion factor $\alpha_{\text{CO}} = 0.8 M_\odot (K \text{ km s}^{-1} \text{ pc}^{-2})^{-1}$ for converting from CO luminosity to molecular gas mass, as estimated for local ultra luminous infrared galaxy (ULIRGs). This conversion factor is very uncertain for high-redshift galaxies. If instead we used the value $\alpha_{\text{CO}} = 4.6 M_\odot (K \text{ km s}^{-1} \text{ pc}^{-2})^{-1}$ which is estimated for local spiral galaxies (Solomon & Vanden Bout 2005), and has also been used for molecular gas estimates in star-forming galaxies at somewhat lower redshifts (e.g. Tacconi et al. 2010), then we would infer six times larger molecular gas masses for the two lensed LBGs from the same CO measurements. Given the small size of the observational sample, and the fact that the model value includes atomic gas, this would make the model roughly compatible with the observational constraints on the molecular gas masses.

We have also investigated the stellar metallicities of LBG descendants. We find that there is a close correlation between mean stellar metallicities of present-day descendants of LBGs and their stellar masses. However, at a given stellar mass, the average metallicity is in most cases almost identical to that of galaxies which are not descended from LBGs in that redshift and UV luminosity range. The only exception to this appears to be for galaxies in the low-mass tail of the descendant stellar mass distribution (for given LBG redshift and UV luminosity range), for which the metallicities of LBG descendants can be higher than for non-descendants. The stellar metallicities of present-day galaxies in the Baugh et al. (2005) model have been studied by Nagashima et al. (2005), who found good agreement between the model predictions and the observed total metallicities of early-type galaxies of different masses.

4.2 Host halo masses

The model predicts that the host halo masses of LBGs and of their present-day descendants show trends with LBG redshift in the same sense as for the stellar masses, i.e. host halo masses for LBGs of a given UV luminosity decrease with increasing redshift, while the host halo masses of their descendants present-day increase. We now discuss this in detail for bright and faint LBGs.

We plot the predicted distributions of host halo masses of LBGs and of their present-day descendants in the left- and right-hand panels of Fig. 7. The upper panels show results for bright LBGs at $z = 3$ and 6, while the lower panels show results for faint LBGs at $z = 3$, 6 and 10.

We see that bright LBGs at $z = 3$ live in haloes with a median mass of $3 \times 10^{11} h^{-1} M_\odot$, while bright LBGs at $z = 6$ reside in...
Figure 6. The stellar mass distribution of LBGs (left) and their $z=0$ descendants (right). The top set of panels show predictions for bright LBGs ($L_{UV} > L_{UV}^*$), and the bottom set of panels faint LBGs ($L_{UV} > 0.1L_{UV}^*$). Results are shown for $z = 3, 6$ and 10 (the last for faint LBGs only) as indicated in each panel. The median stellar mass value is given in the label and is indicated by an arrow. Galaxies are separated into bulge-dominated ($B/T > 0.5$, red histogram) and disc-dominated galaxies ($B/T < 0.5$, blue histogram). The percentage of galaxies in each class is labelled in the same colour as the corresponding histogram. The red and blue distributions add to give the total stellar mass distribution (black histogram).
Figure 7. Host halo mass distribution of LBGs (left) and the host halo mass distribution of their $z=0$ descendants (right). The top set of panels show bright LBGs ($L_{\text{UV}} > L_{\text{UV}}^*$) and the bottom set of panels show faint LBGs ($L_{\text{UV}} > 0.1L_{\text{UV}}^*$). Predictions are shown for $z = 3$, 6 and 10 (the last for faint LBGs only), as indicated in each panel. The median host halo mass value is written in each panel and indicated by an arrow. The red and blue histograms show the predictions for bulge- and disc-dominated galaxies, respectively; the percentage of galaxies in each class is given in the legend, using the same colour.
haloes with a median mass of $2 \times 10^{11} \, h^{-1} \, M_\odot$. The descendants of bright LBGs identified at $z = 3$ are hosted by haloes with a median mass of $1.6 \times 10^{13} \, h^{-1} \, M_\odot$, and the $z = 0$ descendants of bright $z = 6$ LBGs are hosted by haloes with a median mass of $5 \times 10^{13} \, h^{-1} \, M_\odot$.

Faint LBGs at $z = 3$ live in haloes with a median mass of $8 \times 10^{10} \, h^{-1} \, M_\odot$, which is four times larger than the haloes which host faint LBGs at $z = 6$. The distribution of host halo mass for faint LBGs at $z = 10$ is similar to that at $z = 6$. The descendants of faint $z = 3$ LBGs are hosted by haloes with a median mass of $2.5 \times 10^{12} \, h^{-1} \, M_\odot$, and the descendants of faint $z = 6$ LBGs are hosted by haloes with a median mass of $1.3 \times 10^{13} \, h^{-1} \, M_\odot$. Haloes which host the descendants of faint LBGs have a broader distribution in mass than haloes which host the progenitors of bright LBGs.

The host halo masses of LBGs can be observationally constrained through measurements of their clustering. Lacey et al. (2011) show predictions for the linear clustering bias $b$ of LBGs as a function of luminosity and redshift for the same model as used here. For bright LBGs, they find $b \approx 2$ at $z = 3$, increasing to $b \approx 4$ at $z = 6$. As they discuss, these predictions seem broadly consistent with observational estimates of the clustering bias by Adelberger et al. (1998) and Ouchi et al. (2004).

Earlier calculations with a different semi-analytical model suggested that LBGs reside in more massive haloes ($\approx 2 \times 10^{12} \, h^{-1} \, M_\odot$) at $z = 3$ (Baugh et al. 1998; Governato et al. 1998). However, this earlier model had several important differences from the Baugh et al. (2005) model used in the present paper. It did not include the effect of dust extinction on the UV luminosity of the model galaxies, it invoked a much stronger supernova feedback in low circular velocity haloes, it had much less star formation in merger-driven starbursts at high redshift and it did not assume a top-heavy IMF in bursts. The last three of these differences mean that in the Baugh et al. (2005) model, objects can be bright in the UV (and thus be classified as LBGs) while having lower masses than in the Baugh et al. (1998) model.

### 4.3 Morphology

In Section 4.1 we showed that the model predicts a mixture of morphological types (as defined using $B/T$) for LBGs and their descendants. Fig. 8 shows the distribution of $B/T$ for bright and faint LBGs at $z = 3, 6$ and 10, and for their present-day descendants. We see that faint LBGs are typically more bulge dominated in stellar mass than bright LBGs at the same redshift. However, for their descendants, this trend is reversed, with the descendants of bright LBGs being more bulge dominated. There is also a tendency for LBGs at a given luminosity (whether bright or faint) to become more bulge dominated in stellar mass with increasing redshift, and the same trend is also seen for their present-day descendants. The model also predicts that present-day galaxies which are descendants of LBGs in a given redshift and luminosity range have somewhat higher $B/T$ on average than other galaxies of the same stellar mass which are not descended from such LBGs. (We find this result for descendants of both bright and faint LBGs in the redshift ranges $3 < z < 4$ and $6 < z < 7$.)

Observationally, the morphologies of LBGs have mostly been studied by imaging in their rest-frame UV light. Therefore, in Fig. 9, we show the predicted distributions of $B/T$ ratios of LBGs as seen in dust-extincted rest-frame UV luminosity for the same LBG samples at $z = 3, 6$ and 10 as shown in the left-hand panels of Fig. 8. We see from this comparison that the LBGs are typically much more bulge dominated in rest-frame UV light than in stellar mass. The reason for this is that the model LBGs are dominated by starburst systems, for which the young stars producing most of the UV light are assumed to be in a spheroidal component.

![Figure 8](http://mnras.oxfordjournals.org/)

Figure 8. The distribution of $B/T$ for LBGs (left) and their $z = 0$ descendants (right). The distributions for faint LBGs ($L_{UV} > 0.1 L_{UV}^\ast$) are plotted in green and for bright LBGs ($L_{UV} > L_{UV}^\ast$) in black. The top row shows predictions for LBGs identified at $z = 3$ and their descendants. The middle row shows the equivalent predictions for $z = 6$ LBGs and the bottom row for $z = 10$ LBGs. The median $B/T$ value is written in each panel and indicated by an arrow in the same colour as the corresponding histogram.
4.4 Mode of star formation in LBGs

As discussed in Section 2, the model assumes two modes of star formation, quiescent star formation in discs and star formation in bursts triggered by galaxy mergers. Quiescent star formation produces stars with a standard solar neighbourhood IMF, whereas in burst star formation, a top-heavy IMF is adopted. This difference in the choice of IMF between the two modes of star formation has an impact on the abundance and properties of LBGs. In this section, we study correlations between the SFR, stellar mass, halo mass and UV luminosity of galaxies at $z = 3, 6$ and $10$. Our goal is to understand the role of the different star formation channels in the production of LBGs and in determining their nature.

In Fig. 10, we plot all galaxies with stellar masses $M_\star > 10^7 h^{-1} M_\odot$ at $z = 3, 6$ and $10$, showing the relations between stellar mass (plotted on the $x$-axis) and various properties on the $y$-axis: host halo mass (left-hand panel), SFR (middle panel) and far-UV (1500 Å) luminosity in units of $L_{1500}$ (right-hand panel). (We choose a lower mass limit of $M_\star = 10^7 h^{-1} M_\odot$ because the masses of faint LBGs at $z = 10$ extend down to roughly this mass.) The shading in Fig. 10 reflects the square root of the space density of galaxies, becoming darker where more galaxies are found. We gain further insight by dividing the galaxies according to which the mechanism of star formation is dominant (but without the factor 4 used earlier in Section 3.3) and whether the galaxies are central or satellites in their host dark halo. Contours of a given colour enclose 95 per cent of the galaxies passing one of these selection criteria. The red contours delimit the region occupied by 95 per cent of the galaxies in which starbursts are the dominant mode of star formation and blue contours enclose the quiescent population. Similarly, orange contours show where the majority of central galaxies lie and green contours show where the satellite galaxies fall.

Fig. 10 shows that there is a correlation between galaxy stellar mass and host halo mass which is driven by central galaxies. Satellite galaxies in the model, which are mostly quiescent, are responsible for the large scatter in this relation (as suggested by the green contour). There is also a correlation between stellar mass and SFR. In this case the correlation is largely due to galaxies where quiescent star formation dominates (as suggested by the blue contours), with the scatter produced mainly by galaxies which are undergoing bursts of star formation (see also Lagos et al. 2011).

In Fig. 11 we plot similar correlations, but in this case for faint LBGs ($L > 0.1 L_{1500}$) only. Again, we differentiate between galaxies according to the dominant mode of star formation, and make this distinction for the faint and bright LBG samples. Fig. 11 shows that most LBGs are produced by bursts of star formation, as discussed earlier in Section 3.3 and also by Lacey et al. (2011). We again see that the fraction of LBGs dominated by burst star formation increases with redshift. The correlations of halo mass and SFR with stellar mass are seen to be better defined, with less scatter, for LBGs than for samples containing all galaxies. However, the correlation of SFR with stellar mass is also seen to be different (shifted to higher SFRs) for LBGs from that for all galaxies. This is because the LBG samples are dominated by bursts, which have larger SFRs than quiescent galaxies of the same stellar mass.

We have also investigated what fraction of the stellar mass of model LBGs has formed quiescently or in bursts (allowing for the recycling of mass from dying stars). We find that LBGs at $z = 3$ typically formed most of their stellar mass by quiescent star formation, with 57 per cent of faint LBGs and 73 per cent of bright LBGs having formed more than half of their stellar mass quiescently. However, at higher redshifts, the reverse is true: at $z = 6, 93$ per cent
of faint LBGs and 89 per cent of bright LBGs have built at least half of their stellar mass through bursts of star formation after merging with other galaxies, while at $z = 10$, 99 per cent of faint LBGs have built at least half of their stellar mass through bursts. (See also the discussion in Lacey et al. 2011.)

4.5 Comparison with previous work

A key component of our model is the assumption of a non-universal stellar IMF, with a standard IMF adopted in quiescent star formation and a top-heavy IMF in merger-triggered starbursts. This feature of the model precludes any comparison with results presented in terms of stellar masses or SFRs which have been derived assuming a single IMF, as discussed in Section 4.1 above. It also precludes a comparison with most gas dynamics simulations or indeed any semi-analytical models which do not report predictions in terms of the rest-frame UV luminosity. Furthermore, in general, gas dynamics simulations are either not run to the present day or predict present-day mass functions with too many massive galaxies, both of which limit their usefulness for a study of the descendants of LBGs.

The closest analysis to the one presented here is by Guo & White (2009), who used an independent semi-analytical model to examine the properties of LBGs and their descendants. Their study is in a similar spirit to ours, with the parameters of the model set to reproduce the observations of the local galaxy population. The major differences between Guo & White (2009) and our work are the following: (1) the use of a universal IMF in Guo & White; (2) the treatment of dust extinction and (3) the identification of LBGs. The assumption of a universal solar neighbourhood IMF for all modes of star formation in Guo & White (2009) means that the intrinsic UV luminosity will be very different from that predicted by our model for starbursts. The treatment of dust extinction is also very different in the two models. In the Baugh et al. (2005) model, dust extinctions are calculated from a radiative transfer model, with a characteristic optical depth computed for each galaxy based on the size of the disc and bulge, and the mass and metallicity of the cold gas (see Section 2.1). Guo & White (2009) invoke an empirically based dust extinction, with a redshift- and metallicity-dependent term which is tuned to match the observed abundance of LBGs. Their model overestimates the metallicity of the cold gas at $z = 3$ by a factor of 2–3, but agrees better with observations at $z = 0$, and so displays less evolution in gas metallicity with redshift than is observed. Guo & White apply a colour selection to define LBGs in redshift intervals centred on $z = 2.2$ and 3, whereas we have selected samples based on rest-frame UV luminosity. The success of the colour selection in isolating galaxies in a given redshift interval is examined in detail in Gonzalez-Perez et al. (in preparation). At $z = 3$, the colour selection is completely successful at extracting all galaxies brighter than the UV luminosity corresponding to the apparent magnitude limit of a given observation. This means that in this case the UV-selected samples in our work are exactly equivalent to the colour-selected LBG samples in Guo & White. Gonzalez-Perez et al. find that different
colour-selection criteria work with varying degrees of success at higher redshifts.

The apparent magnitude limit applied by Guo & White (2009) corresponds to $L_{UV} \sim 0.4L_{UV}^*$ which is intermediate between our faint and bright samples. The nature of LBGs is very different in the two models. Guo & White find that $z = 3$ LBGs are powered by quiescent star formation and are disc dominated (in terms of stellar mass). Fig. 3 shows that in our model, $z = 3$ LBGs are lit up by starbursts triggered by galaxy mergers. We find that both bright and faint LBGs have fairly flat bulge-to-total stellar mass distributions. Guo & White find that LBGs have higher stellar masses and reside in more massive haloes compared with our model. The descendants of LBGs in Guo & White are 10 times more massive and reside in haloes that are on average 300 times more massive than the host haloes at $z = 3$. We find that the median stellar mass of the descendants of LBGs is $\sim 50$ times larger than the median of the LBGs, and that the present-day host haloes are typically less massive than in the model of Guo & White. Guo & White argue that the descendants are mainly bulge dominated, whereas we find a range of morphologies, with bulge-to-total mass distributions with peaks around 0 and 1. Guo & White do not consider the connection between SMGs and LBGs, which we discuss in Section 5.

5 WHAT IS THE CONNECTION BETWEEN LBGs AND SMGs?

Samples of galaxies at high redshift have been assembled using the Lyman-break technique and through their emission at submillimetre wavelengths. LBGs are selected on the basis of their emission in the rest-frame far-UV. SMGs are detected through their rest-frame far-IR/submillimetre emission, which is produced by the reradiation by dust of light absorbed predominantly in the UV. It is natural to ask what is the overlap, if any, between these two populations at high redshift. A first indication of the answer is suggested by the different redshift ranges of the populations. LBGs selected using ground-based surveys are restricted to have $z \gtrsim 3$. Bright SMGs (with $850$ µm flux $S_{850} > 5$ mJy), on the other hand, have a median redshift of $z \sim 2$ in flux-limited samples. Hence, to study the overlap between the two samples we will be looking at the UV luminosities of galaxies in the high redshift tail of the SMG redshift distribution. Almeida et al. (2010) performed a similar comparison at $z = 2$, and found that $\approx 1$ per cent of galaxies bright in the UV were also predicted to be bright SMGs.

The multiwavelength nature of our semi-analytical model and its self-consistent treatment of dust make it ideally placed to address the overlap between the SMG and LBG populations. Here we examine the predicted relation between $850$ µm flux and UV luminosity at $z = 3$ and $6$. Following the companion to this paper (González et al. 2011), in which we study the evolution of SMGs in the model, we define a galaxy as a bright SMG when its $850$ µm flux exceeds $5$ mJy and as a faint SMG when its $850$ µm flux exceeds $1$ mJy.

The top panels of Fig. 12 show the relation between the $850$ µm flux and the UV luminosity for $z = 3$ and $6$. Following the companion to this paper (González et al. 2011), in which we study the evolution of SMGs in the model, we define a galaxy as a bright SMG when its $850$ µm flux exceeds $5$ mJy and as a faint SMG when its $850$ µm flux exceeds $1$ mJy.

The top panels of Fig. 12 show the relation between the $850$ µm flux and the UV luminosity for galaxies at $z = 3$ and 6, respectively. There is scatter in this relation, particularly at fainter magnitudes. Nevertheless, an increase in the UV luminosity typically corresponds to an increase in the submillimetre flux.

The middle panels of Fig. 12 show the number density of galaxies as a function of their $850$ µm flux, with separate curves for faint and
Figure 12. The relation between the 850 µm flux and the rest-frame far-UV absolute magnitude for galaxies at \(z = 3\) (left-hand panels) and 6 (right-hand panels). Top panels: the median 850 µm flux as a function of the rest-frame far-UV absolute magnitude for galaxies at \(z = 3\). The bars indicate the 10 to 90 percentile range of the distribution in each magnitude bin. The horizontal lines indicate 1 and 5 mJy thresholds, which we used to define faint and bright SMGs in González et al. (2011), and the vertical lines indicate the faint (0.1\(L_{\star\text{UV}}\)) and bright (\(L_{\star\text{UV}}\)) LBG thresholds. Middle panels: the number density of galaxies as a function of the 850 µm flux for faint LBGs (blue histogram) and bright LBGs (red histogram). Bottom panels: the number density of galaxies as a function of the UV absolute magnitude for faint SMGs (\(S_{850} > 1\) mJy, blue histogram) and bright SMGs (\(S_{850} > 5\) mJy, red histogram). In the middle and bottom panels, the arrows indicate the median value of the SMG flux and LBG luminosity, respectively, for each sample. (We omit these arrows where there are too few galaxies in our model sample to reliably calculate the median.)

6 SUMMARY

We have investigated the nature, formation histories and present-day descendants of LBGs in the \(\Lambda\) cold dark matter cosmology that are predicted by the Baugh et al. (2005) version of the GALFORM semi-analytical galaxy formation model. The Baugh et al. model incorporates a top-heavy IMF in starbursts triggered by galaxy mergers, which was found to be necessary to explain the number counts and redshift distribution of SMGs. This model has previously been shown to match the observed far-UV luminosity function of LBGs over the whole range \(z = 3–10\) (Lacey et al. 2011). In this paper, we have investigated a range of other predictions for LBGs from the same galaxy formation model.

We began by investigating the formation histories of LBGs. We found that the majority of LBGs are predicted to be starbursts triggered by minor mergers of galaxies. Using the observed characteristic far-UV luminosity at \(z \approx 3\), \(L_{\star\text{UV}}\), to define samples of model galaxies as bright (\(L_{\star\text{UV}} > L_{\star\text{UV}}\)) or faint (\(L_{\star\text{UV}} > 0.1L_{\star\text{UV}}\)) LBGs, we found that >99 per cent of bright LBGs at \(z = 3–6\) are predicted to be starbursts triggered by minor or major mergers, while this trend is even more pronounced for faint LBGs.
fraction is somewhat lower for faint LBGs (but still > 80 per cent for \( z = 3–10 \)). The importance of starbursts in producing galaxies visible as LBGs is because the far-UV light is dominated by young stars, and in starbursts the SFR is increased over the quiescent value. However, in our model there is also a second important effect that the top-heavy IMF in bursts (as against a solar neighbourhood IMF) further enhances their far-UV luminosities over quiescent galaxies. We then examined the duration of the LBG phase in model galaxies, and found that this is quite short, around 20 Myr for bright LBGs, and somewhat longer, around 20–60 Myr, for faint LBGs. The short time-scales are a result of the LBGs mainly being starbursts.

We next computed the fraction of present-day galaxies that are predicted to have LBG progenitors. We find that 95 per cent of Milky Way mass galaxies had at least one faint LBG progenitor in the redshift range \( 3 < z < 4 \), and 70 per cent at \( 6 < z < 7 \), but this falls to 7 per cent for \( 10 < z < 11 \). For bright LBG progenitors these fractions drop to 40 per cent for LBGs at \( 3 < z < 4 \) and only 3 per cent at \( 6 < z < 7 \). Performing a similar analysis for dark matter haloes, we find that 50 per cent of present-day haloes with masses exceeding \( 4 \times 10^{12} \) or \( 5 \times 10^{13} \) \( M_\odot \) had progenitor haloes hosting bright LBGs at \( 3 < z < 4 \) or \( 6 < z < 7 \), respectively.

We next examined the stellar and halo masses of LBGs, together with their morphologies. The model predicts that bright LBGs at \( z = 3 \) have median stellar mass \( \sim 1 \times 10^{10} \) \( h^{-1} \) \( M_\odot \), falling to \( \sim 5 \times 10^{9} \) at \( z = 6 \), and are typically disc dominated in stellar mass, consistent with their origin in starbursts triggered by minor galaxy mergers. As already noted by Lacey et al. (2011), these stellar masses are lower than the values typically inferred from observational data by fitting stellar population models to broadband photometry, but the observational estimates rely on assuming a solar neighbourhood IMF, while in our model the light from LBGs is dominated by stars formed with a top-heavy IMF during a starburst. The faint LBGs are predicted to have stellar masses roughly an order of magnitude lower than bright LBGs at the same redshift, falling to only \( \sim 1 \times 10^{9} \) \( h^{-1} \) \( M_\odot \) at \( z = 10 \), and to be more bulge dominated than the bright LBGs, especially at the higher redshifts. The haloes hosting bright LBGs at \( z = 3 \) and 6 are predicted to have median masses \( \sim 10^{11} \) \( h^{-1} \) \( M_\odot \), with roughly 10 times lower masses for faint LBGs at the same redshifts.

Following on from this, we then investigated the stellar and halo masses, and morphologies of the present-day descendants of LBGs. The ability to relate different types of high-redshift galaxies to their descendants is one of the unique capabilities of theoretical models of galaxy formation such as ours. We find that the descendants of bright LBGs at \( z = 3 \) have median stellar masses \( \sim 5 \times 10^{10} \) \( h^{-1} \) \( M_\odot \) (i.e. comparable to the Milky Way), and consist of roughly equal numbers of disc- and bulge-dominated systems. They typically inhabit galaxy groups, with halo masses \( \sim 10^{13} \) \( h^{-1} \) \( M_\odot \). Bright LBGs at \( z = 6 \) have descendants which are somewhat more massive \( \sim 1 \times 10^{11} \) \( h^{-1} \) \( M_\odot \) and inhabit somewhat more massive haloes, but are significantly more bulge dominated (\( \sim 75 \) per cent) on average. Faint LBGs are predicted to have present-day descendants which are an order of magnitude lower in stellar and halo mass than the bright LBGs at the same redshift, and contain a larger proportion of disc-dominated systems.

We next looked at the relationship between halo mass, stellar mass, SFR and far-UV luminosity in the model, for all galaxies and for LBGs. We find that LBGs have a tighter relationship between halo mass and stellar mass than galaxies as a whole, due to them being mainly central galaxies. We also find a correlation between SFR and stellar mass that is shifted to higher SFRs compared to that for all galaxies. This results from LBGs being mainly starbursts, which have higher SFRs than quiescent galaxies with the same stellar mass. We also find that although most LBGs in the model are starbursts, their stellar masses are not necessarily dominated by stars formed in bursts. This results in part from the large recycled fraction for stars forming with the top-heavy burst IMF. LBGs at \( z = 3 \) (both faint and bright) are predicted to have formed most of their stellar mass quiescently before the current burst. However, this changes with redshift, so that LBGs at \( z = 6 \) have typically formed most of their stellar mass in bursts (although not necessarily the current burst).

Finally, we used the multiwavelength capabilities of our model to examine the overlap between LBGs and SMGs detected at 850 µm, extending to \( z = 3, 6 \) and the comparison previously carried out by Almeida et al. (2010) at \( z = 2 \). Analogously to the LBGs, we define bright and faint SMGs as having 850 µm fluxes exceeding 5 and 1 mJy, respectively. At \( z \geq 3 \) we are only probing the tail of the redshift distribution of bright SMGs, which peaks at \( z \sim 2 \). We find that only about 1 per cent of bright LBGs at \( z = 3 \) are also predicted to be bright SMGs, and this fraction falls to only 0.1 per cent for bright LBGs at \( z = 6 \), with the median 850 µm flux being about 1 mJy for both redshifts. This explains the difficulty in observationally detecting the submillimetre emission from LBGs (Chapman et al. 2000).

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