S-CURVES IN POLYNOMIAL EXTERNAL FIELDS

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Dedicated to the memory of Andrei Aleksandrovich Gonchar and Herbert Stahl

Abstract. Curves in the complex plane that satisfy the S-property were first introduced by Stahl and they were further studied by Gonchar and Rakhmanov in the 1980s. Rakhmanov recently showed the existence of curves with the S-property in a harmonic external field by means of a max-min variational problem in logarithmic potential theory. This is done in a fairly general setting, which however does not include the important special case of an external field $\text{Re} V$ where $V$ is a polynomial of degree $\geq 2$. In this paper we give a detailed proof of the existence of a curve with the S-property in the external field $\text{Re} V$ within the collection of all curves that connect two or more pre-assigned directions at infinity in which $\text{Re} V \to +\infty$. Our method of proof is very much based on the works of Rakhmanov on the max-min variational problem and of Martínez-Finkelshtein and Rakhmanov on critical measures.

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1. Introduction

The concept of $S$-property for logarithmic potentials appeared for the first time in the works of H. Stahl [39–42], in the study of the limiting distribution of poles of Padé approximants, see also [43] for a recent survey and extension of these results.

Inspired by Stahl’s works, Gonchar and Rakhmanov extended the $S$-property to situations with external field [21]. They then characterized the limiting distribution of zeros of certain non-hermitian orthogonal polynomials, subject to the existence of a certain curve with the $S$-property in an external field - the $S$-curve, see Section 2.2 below for a precise definition.

Very recently, Rakhmanov returned to the question of existence of $S$-curves from a general perspective [33]. He considered an associated max-min equilibrium problem in logarithmic potential theory, and in an “ideal” situation he proved that this max-min problem has a solution and the solution has the $S$-property. This general approach by Rakhmanov is very similar in spirit to his previous work with Kamvissis [25], where a max-min problem for Green energy in a particular external field was considered.

In [33], Rakhmanov also pointed out some examples where his proposed “ideal” situation does not apply, although similar considerations should also lead to the existence of a curve with the $S$-property. The present work is aimed at studying this problem in a very particular, but important, model, namely when the external field is given by the real part of a polynomial.

Following the approach of Rakhmanov, for a polynomial $V$ we consider a max-min problem of the form

$$\max_{\Gamma} \min_{\mu \subset \Gamma} \left[ \int \int \log \frac{1}{|x-y|} d\mu(x)d\mu(y) + \int \text{Re} V(x) d\mu(x) \right],$$

where the maximum is taken over a suitable class of contours $\Gamma$ (see Section 2.3 below for precise definitions), and the infimum is taken over the set of Borel probability measures $\mu$ on $\Gamma$. The aim of this work is to prove that the max-min problem has a solution, and this solution leads to an $S$-curve $\Gamma$.

The max-min approach is not the only possible approach to the existence of the $S$-curve. In some particular cases, it is possible to construct the $S$-curve explicitly from specific properties of the problem at hand, see e.g. [1, 3, 7, 11, 14, 16] for recent contributions. In general, the explicit determination of the $S$-curve - or more specifically the determination of the support of its equilibrium measure - is a hard problem, and some numerical studies have also been carried out, see [11].

Motivated by questions in Random Matrix Theory, Bertola [9] also studied the existence of the $S$-curve for polynomial external field in the same framework as the present paper. By using deformation techniques, Bertola was able to show the existence of the $S$-curve. However, his approach just works when the underlying equilibrium measure vanishes as a square root on the endpoints of its support, although it is very likely that his approach could also be adapted to remove this restriction. It is worth emphasizing that in our work there is no restriction on the underlying equilibrium measure - so we get the same result as Bertola but valid in the general situation.

An extension of the concept of $S$-curves to vector equilibrium problems appears in the context of Hermite-Padé approximation, see e.g. [5, 32, 33]. Vector equilibrium problems also appear in certain problems in random matrix theory, for example for
random matrix models with external source \[4, 12\] and coupled random matrices \[19\], see \[26\] for an overview. The asymptotic analysis has been mostly restricted to cases with enough symmetry so that the relevant contours are contained in the real or imaginary axis. A suitable existence theory for curves with S-property appropriate to vector equilibrium problems with external fields, and possibly also upper constraints, would be a major step towards the analysis of situations with less symmetry. See \[6\] for a first example in this direction.

2. Background and statement of the main result

To state our main result, we first need to introduce a few notations and notions. Throughout this paper, \(V\) always denotes a complex polynomial of degree \(N \geq 2\) and \(\varphi = \text{Re} V\).

Also, we denote
\[
D_R(z_0) = \{ z \in \mathbb{C} | |z - z_0| < R \}, \quad D_R = D_R(0).
\]

2.1. Notions from potential theory. Given a finite Borel measure \(\mu\) on \(\mathbb{C}\), its logarithmic potential at \(x \in \mathbb{C}\) is defined by
\[
U_\mu(x) = \int \log \frac{1}{|x - y|} \, d\mu(y),
\]
and its logarithmic energy by
\[
I(\mu) = \iint \log \frac{1}{|x - y|} \, d\mu(x) d\mu(y) = \int U_\mu(x) \, d\mu(x),
\]
whenever these integrals make sense.

For a closed set \(F \subset \mathbb{C}\) and \(\varphi\) the real part of a polynomial as before, we define \(M_1(F)\) as the set of Borel probability measures \(\mu\) supported on \(F\), satisfying the growth condition
\[
\int \log(1 + |x|) \, d\mu(x) < +\infty,
\]
and by \(M_1^\varphi(F)\) we denote the subset of \(M_1(F)\) consisting of all measures \(\mu\) for which \(I(\mu)\) is finite and \(\varphi \in L^1(\mu)\). The condition \(\text{(2.1)}\) assures the quantities \(U_\mu(x)\) and \(I(\mu)\) are well defined elements of \((-\infty, +\infty]\) and \(U_\mu(z)\) is finite for a.e. \(z \in \mathbb{C}\) with respect to planar Lebesgue measure.

For \(\mu \in M_1^\varphi(F)\), the quantity
\[
I^\varphi(\mu) = I(\mu) + \int \varphi(x) \, d\mu(x),
\]
is called the weighted logarithmic energy of \(\mu\) in the external field \(\varphi\), or just weighted energy of \(\mu\).

The following minimization problem is classical \[36\]: find a measure \(\mu^\varphi = \mu^\varphi, F \in M_1^\varphi(F)\) for which the infimum
\[
I^\varphi(F) = \inf_{\mu \in M_1^\varphi(F)} I^\varphi(\mu)
\]
is attained, i.e., for which \(I^\varphi(\mu^\varphi) = I^\varphi(F)\). The measure \(\mu^\varphi\) is called the equilibrium measure of the set \(F\) in the external field \(\varphi\). For the case without external field, that is, when \(\varphi \equiv 0\), the measure \(\mu^0\) is also called the Robin measure of \(F\), the
value $I(F)$ is the Robin constant of $F$, the constant $e^{-I(F)}$ is called (logarithmic) capacity of $F$ and denoted by $\text{cap} F$.

For $\varphi$ the real part of a polynomial as before, if $\Gamma$ is a contour in $\mathbb{C}$ satisfying the growth condition
\[ \lim_{z \to \infty, z \in \Gamma} \varphi(z) = +\infty, \] (2.3)
the equilibrium measure $\mu_{\varphi, \Gamma}$ uniquely exists and can be characterized through the Euler-Lagrange variational conditions. It is the only measure in $\mathcal{M}^E_{\varphi}(\Gamma)$ for which there exists a constant $l$ satisfying
\[ U^\mu(z) + \frac{1}{2}\varphi(z) = l, \quad z \in \text{supp} \mu, \] (2.4)
\[ U^\mu(z) + \frac{1}{2}\varphi(z) \geq l, \quad z \in \Gamma. \] (2.5)

2.2. S-property.

Definition 2.1. A contour $\Gamma$ satisfying the growth condition (2.3) is said to possess the S-property in the external field $\varphi$ if its equilibrium measure $\mu_{\varphi, \Gamma}$ in the external field satisfies the following. There is a finite set $E$ such that $\text{supp} \mu \setminus E$ is locally an analytic arc at each of its points, and its potential $U^\mu = U^\mu_{\varphi, \Gamma}$ satisfies
\[ \frac{\partial}{\partial n_+} \left( U^\mu + \frac{1}{2}\varphi \right)(z) = \frac{\partial}{\partial n_-} \left( U^\mu + \frac{1}{2}\varphi \right)(z), \quad z \in \text{supp} \mu \setminus E, \] (2.6)
where $n_+$ are the unit normal vectors to $\text{supp} \mu$ at $z$.

If a contour has the S-property in the external field $\varphi$, we also call it an S-curve in the external field $\varphi$. If the external field $\varphi$ is understood (as it is in this paper), we briefly call it an S-curve.

As it is clear already from the Definition 2.1, the S-property is intrinsic to the equilibrium measure in the external field of the contour. If we modify the contour away from the support, keeping the equilibrium measure in the external field the same for the modified contour, the S-property is preserved.

Although in the present work we are just interested in the S-property stated in terms of logarithmic potentials, its analogue for potentials coming from other kernels, such as the Green potential, have also their own interest [25,30,34].

2.3. The class of admissible contours. There exist $N$ sectors $S_1, \ldots, S_N$ in the complex plane along which $\varphi(z) = \text{Re} V(z) \to +\infty$ as $z \to \infty$ in $S_j$. These sectors are determined by the leading coefficient of $V$. If $V(z) = a_0 z^N + \cdots, a_0 \neq 0$, then
\[ S_j = \left\{ z \in \mathbb{C} \mid \left| \arg z - \theta_j \right| < \frac{\pi}{2N} \right\}, \quad \theta_j = -\frac{\arg a_0}{N} + \frac{2\pi(j-1)}{N}, \] (2.7)
for $j = 1, \ldots, N$.

We say that a set $F \subset \mathbb{C}$ stretches out to infinity in the sector $S_j$ if there exist $\epsilon > 0$ and $r_0 > 0$ such that for every $r > r_0$ there is $z \in F$ with
\[ |z| = r, \quad \left| \arg z - \theta_j \right| < \frac{\pi}{2N} - \epsilon. \]
Thus in particular we have $\varphi(z) \to +\infty$ if $z \to \infty$ within $S_j \cap F$.

A partition $\mathcal{P}$ of $\{1, \ldots, N\}$ is called noncrossing if it satisfies the following property:
\[ j \overset{\mathcal{P}}{\sim} j' \land k \overset{\mathcal{P}}{\sim} k' \land j < k < j' < k' \implies j \overset{\mathcal{P}}{\sim} k. \]
Here $\sim$ denotes the equivalence relation that is associated with the partition $\mathcal{P}$.

The term noncrossing becomes more evident in a graphical representation of a partition, see Figure 1.

![Figure 1](https://example.com/figure1.png)

**Figure 1.** Graphical representation of the crossing partition $\mathcal{P} = \{\{1,4,5\}, \{2,3,6\}\}$ (on the left) and of the noncrossing partition $\mathcal{P} = \{\{1,5,6\}, \{2,3\}, \{4\}\}$ (on the right).

To any noncrossing partition we associate a collection $\mathcal{T}$ of admissible contours.

**Definition 2.2.** Let $\mathcal{P}$ be a noncrossing partition of $\{1,\ldots,N\}$ and $\mathcal{P}_0$ the subset of $\mathcal{P}$ obtained by removing from $\mathcal{P}$ all singleton sets. We assume $\mathcal{P}_0 \neq \emptyset$. We associate with $\mathcal{P}$ the collection $\mathcal{T}(\mathcal{P})$ of admissible contours $\Gamma$ defined as follows.

i) Each $\Gamma \in \mathcal{T}(\mathcal{P})$ is a finite union of $C^1$ Jordan arcs.

ii) Each $\Gamma \in \mathcal{T}(\mathcal{P})$ has at most $|\mathcal{P}_0|$ connected components, all of them unbounded and stretching out to infinity in at least two of the sectors $S_1, \ldots, S_N$.

iii) For each $A \in \mathcal{P}_0$ there is a connected component $\Gamma_A$ of $\Gamma$ that stretches out to infinity in each sector $S_j$ with $j \in A$.

iv) If $A \in \mathcal{P} \setminus \mathcal{P}_0$ then there exists $R > 0$ for which

$$\Gamma \cap (S_j \setminus D_R) = \emptyset, \quad j \in A.$$ 

For simplicity, we will mostly write $\mathcal{T}$ instead of $\mathcal{T}(\mathcal{P})$.

Figure 2 shows all noncrossing partitions (up to rotation) for a polynomial of degree $N = 5$ and representative contours $\Gamma$ from each of the collections $\mathcal{T}(\mathcal{P})$.

The restriction to noncrossing partitions in Definition 2.2 is not essential. We could as well have defined the admissible contours for arbitrary partitions, but then the class of admissible contours obtained for a given crossing partition would coincide with the class of admissible contours for some noncrossing partition.

2.4. **The main theorem.** For the polynomial $V$ and a fixed noncrossing partition $\mathcal{P}$ as above, the *max-min energy problem* for the pair $(V, \mathcal{T}(\mathcal{P}))$ asks for finding a contour $\Gamma_0 \in \mathcal{T}$ for which $I^x(\cdot)$, see (2.2), attains its maximum on $\mathcal{T}$. That is,

$$I^x(\Gamma_0) = \max_{\Gamma \in \mathcal{T}} I^x(\Gamma) = \max_{\Gamma \in \mathcal{T}} \min_{\mu \in \mathcal{M}_1^x(\Gamma)} I^x(\mu). \quad (2.8)$$

In what follows we use

$$C^\mu(z) = \int \frac{d\mu(x)}{x-z} \quad (2.9)$$

to denote the Cauchy Transform of a measure $\mu$. The main result of this paper is the following.

**Theorem 2.3.** Let $V$ be a polynomial of degree $N \geq 2$ and let $\mathcal{P}$ be a noncrossing partition with class $\mathcal{T}(\mathcal{P})$ of admissible contours as in Definition 2.2. Then there exists a solution $\Gamma_0 \in \mathcal{T}(\mathcal{P})$ to the max-min energy problem (2.8) for $(V, \mathcal{T}(\mathcal{P}))$.

The contour $\Gamma_0$ has the $S$-property in the external field $\varphi = \text{Re}V$. 

The equilibrium measure \( \mu_0 := \mu^{\varphi, \Gamma_0} \) of \( \Gamma_0 \) in the external field \( \varphi \) is supported on a finite union of analytic arcs that are critical trajectories of the quadratic differential \(-R(z)dz^2\), where

\[
R(z) = \left( C^{\mu_0}(z) + \frac{V'(z)}{2} \right)^2, \quad z \in \mathbb{C} \setminus \text{supp} \mu_0, \tag{2.10}
\]

is a polynomial of degree \( 2N - 2 \).

Furthermore \( \mu_0 \) is absolutely continuous with respect to arclength and

\[
d\mu_0(s) = \frac{1}{\pi i} R_+(s)^{1/2} ds. \tag{2.11}
\]

See Section 2.5 for a brief discussion on quadratic differentials and their critical trajectories.
In (2.11) we use a complex line element $ds$ on $\Gamma_0$, which induces an orientation on $\Gamma_0$. Then $R_+(s)^{1/2}$ denotes the limiting value of

$$R(z)^{1/2} = \int \frac{d\mu_0(x)}{x-z} + \frac{V'(z)}{2}$$

as $z \to s \in \Gamma_0$ from the left-hand side.

The contour $\Gamma_0$ that is the maximizer for the max-min energy property in Theorem 2.3 is not unique. We can modify the contour $\Gamma_0$ outside the support of $\mu_0$ slightly, preserving the equilibrium measure in the external field. In this way we obtain a new contour that also solves the max-min problem for $(V, T(\mathcal{P}))$.

Although the contour is not unique, the equilibrium measure in the external field $\mu_0$ and its support turn out to be uniquely determined by $V$ and $\mathcal{P}$. We do not know of a direct way to prove this, but there is an indirect way using orthogonal polynomials. Assume for simplicity that $\mathcal{P}_0 = \{\{i,j\}\}$ for some $1 \leq i < j \leq N$.

This means that any contour $\Gamma \in T(\mathcal{P})$ connects two fixed distinguished sectors at infinity where $\text{Re} V(z) \to +\infty$.

The (non hermitian) bilinear form defined on polynomials $p$ and $q$,

$$\langle p, q \rangle = \langle p, q \rangle_{V,n} = \int_{\Gamma} p(z)q(z)e^{-nV(z)}dz$$

is then well defined and independent of the choice of contours $\Gamma \in T(\mathcal{P})$. A polynomial $p_n$ of degree at most $n$ is orthogonal with respect to $\langle \cdot, \cdot \rangle$ if

$$\langle p_n, q \rangle = 0,$$

for any polynomial $q$ of degree $\leq n - 1$. Associated to the polynomial $p_n$ is its normalized zero counting measure $\chi_n$

$$\chi_n = \frac{1}{\deg p_n} \sum_{p_n(z)=0} \delta_z.$$

In studying $(\chi_n)$, Gonchar and Rakhmanov proved

**Theorem 2.4** ([21, Theorem 3]). *If there exists $\Gamma \in T(\mathcal{P})$ with the S-property in the external field $\varphi = \text{Re} V$ and $\mathbb{C} \setminus \text{supp} \mu^\varphi, \Gamma$ is connected, then*

$$\chi_n \xrightarrow{s} \mu^\varphi, \Gamma.$$

In other words, if a contour $\Gamma \in T(\mathcal{P})$ has the S-property, then its equilibrium measure in the external field is the weak limit of $(\chi_n)$.

By Theorem 2.3, any contour $\Gamma \in T(\mathcal{P})$ with the max-min property has the S-property. Also $\mathbb{C} \setminus \text{supp} \mu^\varphi, \Gamma$ is connected, since otherwise we could remove self-intersecting subarcs of $\Gamma$ and increase its energy. Thus, if $\Gamma_0, \Gamma_1$ are two solutions to the max-min problem $(V, T(\mathcal{P}))$, then by Theorem 2.4 both equilibrium measures $\mu^\varphi, \Gamma_0$ and $\mu^\varphi, \Gamma_1$ are the weak limit of the sequence $(\chi_n)$. By uniqueness of limits, it follows that $\mu^\varphi, \Gamma_0 = \mu^\varphi, \Gamma_1$.

The above argument applies if $\mathcal{P}_0 = \{\{i,j\}\}$ for some $1 \leq i < j \leq N$. For other choices of $\mathcal{P}_0$, we need to modify the bilinear form $\langle \cdot, \cdot \rangle$ to get a right family of orthogonal polynomials $(p_n)$. For more details, we refer to [9][10].
2.5. Quadratic differentials. A polynomial $R$ of degree $2N - 2$, $N \geq 2$, defines on the Riemann sphere $\mathbb{C}$ the quadratic differential $-R(z)dz^2$. A basic reference for quadratic differentials is the book by Strebel [44].

The critical points of $-R(z)dz^2$ are the zeros of $R$, whose orders as critical points are the same as their respective orders as zeros of $R$, and the point $z = \infty$, which is a pole of order $2N + 2$. The other points in $\mathbb{C}$ are regular points of $-R(z)dz^2$.

An arc $\gamma$ is called a horizontal arc of $-R(z)dz^2$ if, along $\gamma$

$$-R(z)dz^2 > 0.$$ 

At a regular point, this condition can be locally stated as

$$\text{Re} \int_{z_0}^{z} \sqrt{R(s)} \, ds \equiv \text{const}, \quad z \in \gamma.$$ 

A maximal horizontal arc is called a horizontal trajectory, or shortly trajectory, of $-R(z)dz^2$, and a trajectory that ends at a zero of $R$ is called a critical trajectory.

Also of interest are the vertical trajectories, which are defined instead by the requirement

$$R(z)dz^2 > 0,$$

which locally at a regular point is equivalent to

$$\text{Im} \int_{z_0}^{z} \sqrt{R(s)} \, ds \equiv \text{const}.$$ 

The local structure of trajectories is well understood. Through any regular point $z_0$ passes exactly one horizontal and one vertical trajectory, which are analytic arcs that intersect at $z_0$ only, where they are orthogonal.

From a zero of $-R(z)dz^2$ of order $p$ emanate exactly $p + 2$ trajectories, at equal consecutive angles $\frac{2\pi}{p+2}$.

At the pole $z = \infty$ of order $2N + 2$, the local behavior of trajectories is as follows [44, Theorem 7.4]. There are $2N$ directions at $\infty$ forming equal angles, such that any trajectory ending at $\infty$ does so in one of these directions. There is a neighborhood $U$ of $\infty$ such that any trajectory entering $U$ ends at $\infty$ in at least one direction, and trajectories entirely contained in $U$ end at $\infty$ in two consecutive directions, see Figure 3, which shows the trajectories at infinity for a pole of order 10 after inversion $z \mapsto \frac{1}{z}$.

2.6. Rakhmanov’s results and overview of the rest of the paper. As mentioned in the Introduction, our work is very much influenced by the recent work of Rakhmanov [33], which we now explain.

The first result obtained by Rakhmanov in [33], which is relevant to us is the following.

**Theorem 2.5 ( [33, Theorem 9.2]).** Suppose $F$ is closed in $\mathbb{C}$. If

$$-\infty < I^\phi(F) < +\infty,$$

then the equilibrium measure $\mu^\phi,F$ in external field $\phi$ exists and is unique.

In contrast with previous results, Theorem 2.5 does not contain any assumption on the growth of $\phi$ at $\infty$. Actually the theorem in [33] is more general as it also deals with situations where the external field is continuous except at a finite number of points, and no growth restriction near these singularities.
Rakhmanov considers the energy functional

$$F \mapsto I^\varphi(F)$$

defined on a metric space of compacts, which we explain next.

The Riemann sphere $\mathbb{C}$ is mapped to the sphere $\mathcal{S} \subset \mathbb{R}^3$ centered at $(0,0,1/2)$ with radius $\frac{1}{2}$,

$$\mathcal{S} = \left\{(x_1, x_2, x_3) \in \mathbb{R}^3 \mid x_1^2 + x_2^2 + (x_3 - \frac{1}{2})^2 = \frac{1}{4}\right\},$$

via the inverse stereographic map $L: \mathbb{C} \to \mathcal{S}$ given by $L(\infty) = (0,0,1)$ and

$$L(z) = \left(\frac{\text{Re} z}{1+|z|^2}, \frac{\text{Im} z}{1+|z|^2}, \frac{|z|^2}{1+|z|^2}\right), \quad z \in \mathbb{C}.$$  

The hyperbolic metric on $\mathbb{C}$ is then given by

$$d(z, w) = ||L(z) - L(w)||,$$

where the norm $||\cdot||$ on the right-hand side above is the Euclidean norm in $\mathbb{R}^3$. The following relation holds

$$d(z, w) = \frac{|z-w|}{\sqrt{1+|z|^2}\sqrt{1+|w|^2}}, \quad z, w \in \mathbb{C}.$$

This metric induces a metric on compact subsets of $\mathbb{C}$. If $K_1, K_2 \subset \mathbb{C}$ are compacts, the Hausdorff metric between them is given by the formula

$$d_H(K_1, K_2) = \max \left\{ \sup_{x \in K_1} \text{dist}(x, K_2), \sup_{y \in K_2} \text{dist}(y, K_1) \right\}, \quad (2.12)$$

where the distance between points and sets above is the usual one induced by $d$,

$$\text{dist}(x, K) = \inf_{y \in K} d(x, y).$$
Alternatively, denoting 

\[ \delta(K) = \{ x \in \overline{C} \mid \text{dist}(x, K) < \delta \}, \]

it can be equivalently defined by

\[ d_H(K_1, K_2) = \inf \{ \delta > 0 \mid K_1 \subset (K_2)_\delta, K_2 \subset (K_1)_\delta \}. \]

The Hausdorff metric is thus defined on closed sets of \( \overline{C} \). Naturally it induces a metric on closed sets of \( C \) by adding the point at infinity to unbounded closed sets. In what follows, when we talk about the Hausdorff metric on closed subsets of \( \overline{C} \), we mean the metric obtained with this natural identification to closed subsets of \( C \).

We equip \( T \) with the Hausdorff metric (2.12) induced by the hyperbolic distance. So whenever we refer to a topological notion on \( T \) such as convergence or continuity, we will always mean with respect to the Hausdorff metric just introduced.

Suppose \( F \) is a class of closed subsets of \( C \) for which there exists an absolute constant \( c \) such that every compact set in \( F \) has at most \( c \) connected components. Equip \( F \) with the Hausdorff metric just introduced. For such classes \( F \), Rakhmanov proves the following.

**Theorem 2.6** ([33, Theorem 3.2]). *The energy functional* \( \mathcal{I} : \mathcal{F} \to [-\infty, +\infty] \) *is upper semicontinuous.*

If the class \( F \) is closed, then it follows from (2.6) there exists a set \( F_0 \in F \) maximizing \( \mathcal{I}^{\varphi} \). However, it is easy to see that the class of admissible contours \( \mathcal{T} = \mathcal{T}(\mathcal{P}) \) introduced in Definition 2.2 is not closed in the Hausdorff metric. An easy way to see this is through space-filling curves. One can approximate sets of large planar measure with smooth contours, and the limiting set is not a smooth contour anymore.

A natural idea would then be to work with the closure \( \overline{T} \) of the class \( T \) in the Hausdorff metric. There is however an issue with that, since sets in \( \overline{T} \) may stretch out to infinity in unlikely directions, and we might then recover the S-curve for the admissible class \( \mathcal{T}(\mathcal{P}') \) corresponding to another choice of partition \( \mathcal{P}' \).

Our technique to overcome this last issue is to introduce a suitable subclass \( \mathcal{T}_M \subset \mathcal{T} \), see (5.4), consisting of contours that do not enter a “forbidden region” \( \Delta_M \). Then we take the closure \( \mathcal{F}_M = \mathcal{T}_M \), see also (5.5). For the closed class \( \mathcal{F}_M \), Theorem 2.6 assures us that there exists a set \( F_0 \) maximizing the energy functional \( \mathcal{I}^{\varphi}(\cdot) \) on \( \mathcal{F}_M \). Moreover, with \( M \) chosen large enough, we can guarantee that the set \( F_0 \) does not come to the boundary of the forbidden region \( \Delta_M \), which implies that small deformations of \( F_0 \) in any direction also belong to \( \mathcal{F}_M \).

Any \( C^2 \) function \( h : \mathbb{C} \to \mathbb{C} \) generates a one-parameter deformation \( F^t_0 \) of this maximizer set by

\[ F^t_0 = \{ z + th(z) \mid z \in F \}, \quad t \in \mathbb{R}. \tag{2.13} \]

Now, since \( F^t_0 \) also belongs to \( \mathcal{F} \) for \( t \) sufficiently close to 0, we have that (the limit actually exists)

\[ \lim_{t \to 0} \frac{\mathcal{I}^{\varphi}(F^t_0) - \mathcal{I}^{\varphi}(F_0)}{t} = 0, \]

because \( F_0 \) is a maximizer set. Rakhmanov [33] then proves that this limit implies that the equilibrium measure \( \mu^{\varphi,F_0} \) in the external field \( \varphi \) is critical, in the sense introduced by Martínez-Finkelshtein and Rakhmanov [29], see also Section 3 below.

As it was already observed by Martínez-Finkelshtein and Rakhmanov [29], the Cauchy transform of a critical measure is an algebraic function. In the present
context, this means that the Cauchy transform \( C^{\mu_0} \) of the equilibrium measure in external field satisfies an algebraic equation of the form:

\[
\left( C^{\mu_0}(z) + \frac{V'(z)}{2} \right)^2 = R(z), \quad \text{for a.e. } z \in \mathbb{C},
\]

where \( R \) is a polynomial of degree \( 2n - 2 \), see Proposition 3.7 below. The main technical difference between [29] and the present work is that in [29] this is obtained under the assumption that the critical measure is compactly supported, while here we obtain the algebraic equation first, and as one of its consequences we find that \( \mu_0 \) must be compactly supported. As in [29], the algebraic equation also implies that the measure \( \mu_0 \) is supported on a finite union of analytic arcs and its density can be recovered from the polynomial \( R \), see Propositions 3.8, 3.9. Also, the S-property ultimately follows from the algebraic equation, see Corollary 3.11.

As a final step, once some properties of the equilibrium measure in external field \( \mu_0 \) of the maximizer set \( F_0 \) are known, we drop the parts of \( F_0 \) that do not belong to \( \text{supp} \mu_0 \), and we give in Section 5.3 a construction of a curve \( \Gamma_0 \in \mathcal{T}(P) \) whose equilibrium measure in the external field \( \varphi \) is the desired measure \( \mu_0 \).

The rest of the paper is organized as follows. In Section 3 we discuss critical measures. The results therein presented are similar to the ones obtained for the first time by Martínez-Finkelshtein and Rakhmanov [29]. However, as mentioned before and further explained later, the setup in [29] is not the same as ours, so we preferred to present detailed proofs in Section 3. In Section 4 we explain the link between critical measures and critical sets, following the lines of [33]. Section 5 is devoted to the proof of Theorem 2.3.

### 3. Critical Measures

As before, \( V \) is a polynomial of degree \( N \geq 2 \) and the external field \( \varphi \) is given by

\[
\varphi(z) = \text{Re} V(z), \quad z \in \mathbb{C}.
\]

With proper modifications, all the results and proofs in this section are also valid for external fields given by the real part of an analytic function, possibly multivalued, with rational (and so single-valued) derivative, but for our purposes it is enough to consider the polynomial case.

The results we are going to show here were proven before by Martínez-Finkelshtein and Rakhmanov [29] under the assumption that the measures involved are a priori compactly supported and the external field has a rational derivative with simple poles. The techniques we are using are also similar, but sometimes simplified and adapted to our needs. The main difference here is that we do not impose a priori that a critical measure must have compact support. Ultimately, we do find that critical measures must be compactly supported (for external fields as in (3.1)), and then the results here are the same as the ones obtained in [29].

#### 3.1. Derivative of the energy functional

Consider the set of test functions \( C_\mathbb{C}^2 \), consisting of \( C^2 \) complex-valued functions in \( \mathbb{C} \) with compact support. For a function \( h \in C_\mathbb{C}^2 \) and \( t \in \mathbb{R} \) denote

\[
h_t(z) = z + th(z), \quad z \in \mathbb{C}.
\]
Any function $h \in C^2$ generates local variations of sets $F \mapsto F^t = h_t(F)$ and also of measures $\mu \mapsto \mu^t$, where $\mu^t$ is the pushforward measure induced by $h_t$. If $\mu$ has finite weighted logarithmic energy then the same is true for $\mu^t$.

For a measure $\mu \in M^2(\mathbb{C})$ and $h \in C^2$, we define the derivative of the weighted energy functional at $\mu$ in the direction of $h$ by

$$D_h I^\varphi(\mu) = \lim_{t \to 0} \frac{I^\varphi(\mu^t) - I^\varphi(\mu)}{t},$$

whenever the limit exists. A first result is that $D_h I^\varphi(\mu)$ indeed always exists.

**Proposition 3.1.** Given $\mu \in M^2(\mathbb{C})$ and $h \in C^2$, the derivative $D_h I^\varphi(\mu)$ exists and is given by

$$D_h I^\varphi(\mu) = -\text{Re} \left( \iint \frac{h(x) - h(y)}{x - y} \ d\mu(x)d\mu(y) - \int V'(x) h(x) \ d\mu(x) \right).$$

**Proof.** The proof here is the same as in [29, Proof of Lemma 3.7]. By the Change of Variables Formula

$$I^\varphi(\mu^t) = \iint \log \left| \frac{x - y + t(h(x) - h(y))}{x - y} \right| \ d\mu(x)d\mu(y) + \int \varphi(x + th(x)) \ d\mu(x),$$

which implies

$$I^\varphi(\mu^t) - I^\varphi(\mu) = -\iint \log \left| 1 + t \frac{h(x) - h(y)}{x - y} \right| \ d\mu(x)d\mu(y) + \int (\varphi(x + th(x)) - \varphi(x)) \ d\mu(x).$$

Since $h \in C^2$, we have for $t \to 0$,

$$\log \left| 1 + t \frac{h(x) - h(y)}{x - y} \right| = \text{Re} \left( \log \left( 1 + t \frac{h(x) - h(y)}{x - y} \right) \right) = t \text{Re} \left( \frac{h(x) - h(y)}{x - y} \right) + O(t^2),$$

with the implicit constant in $O(t^2)$ uniform in $x, y \in \mathbb{C}$, and analogously

$$V(x + th(x)) - V(x) = tV'(x) h(x) + O(t^2),$$

where again the $O$ term is uniform in $x \in \mathbb{C}$. This implies

$$\varphi(x + th(x)) - \varphi(x) = t \text{Re} \left( h(x) V'(x) \right) + O(t^2) \quad \text{as } t \to 0,$$

and the result follows by plugging this last equation and (3.6) into (3.5).  

\[ \Box \]

### 3.2. Critical measures.

**Definition 3.2.** We say a measure $\mu \in M^2(\mathbb{C})$ is $\varphi$-critical if

$$D_h I^\varphi(\mu) = 0$$

for every $h \in C^2$.

The definition of $\varphi$-critical measures introduced in [29] allows the external field to have singularities. In that situation, the test functions $h$ appearing in Definition 3.2 should also vanish on a set $A$ of zero capacity containing the singular points of $\varphi$, leading to the notion of $(\varphi, A)$-critical measures. In [29] it is also imposed that critical measures must have compact support. In the case considered here the only singularity of the external field is at $z = \infty$, and for convenience we are
initially considering functions that are not just vanishing at infinity but also in a neighborhood of it, that is, compactly supported functions.

Considering \( h \) and \( \text{i}h \) separately, we easily obtain the following from Proposition 3.1.

**Corollary 3.3.** A measure \( \mu \in \mathcal{M}_1^\varphi(\mathbb{C}) \) is \( \varphi \)-critical if, and only if,
\[
\int \int \frac{h(x) - h(y)}{x - y} \, d\mu(x) \, d\mu(y) = \int V'(x) h(x) \, d\mu(x) \tag{3.7}
\]
for every \( h \in C^2 \).

Equation (3.7) appears frequently in the context of random matrices, although usually in different forms and names as for example Loop equations, Schwinger-Dyson equations or Ward identities, see [2,18,20].

We need to extend (3.7) to larger classes of test functions.

**Lemma 3.4.** Let \( \mu \) be a \( \varphi \)-critical measure. Then equation (3.7) remains valid for \( h \in C^2 \) satisfying
\[
h(z) = O\left(\frac{1}{z^{N-1}}\right), \quad z \to \infty. \tag{3.8}
\]

**Proof.** Introduce the function \( \Theta_n : [0, +\infty) \to \mathbb{R} \) by
\[
\Theta_n(t) = \begin{cases} 1, & t \leq n, \\ \frac{1}{32n^3} (3n-t)^3 (2n^2 - 3nt + 3t^2), & n < t < 3n, \\ 0, & t \geq 3n. \end{cases}
\]
Then \( \Theta_n \in C^2 \), \( \Theta_n(t) \xrightarrow{n \to \infty} 1 \) and
\[
0 \leq \Theta_n(t) \leq 1, \quad |\Theta_n'(t)| < \frac{1}{n}, \quad t \geq 0. \tag{3.9}
\]

For \( h \in C^2 \) satisfying (3.8), define
\[
h_n(x) = \Theta_n(|x|) h(x), \quad x \in \mathbb{C}.
\]
Then \( h_n \) belongs to \( C^2 \), so that equation (3.7) is valid for it. From assumption (3.8) and the fact that \( \deg V = N \), we see
\[
V'(x) h(x) = \mathcal{O}(1), \quad x \to \infty,
\]
and from the definition of \( h_n \) and the Dominated Convergence Theorem, it is easily seen
\[
\int V'(x) h_n(x) \, d\mu(x) \xrightarrow{n \to \infty} \int V'(x) h(x) \, d\mu(x). \tag{3.10}
\]

The measure \( \mu \) has finite logarithmic energy, so it has no mass points. In particular, the diagonal \( \{ (x,x) \mid x \in \mathbb{C} \} \) has zero \( \mu \times \mu \) measure, implying
\[
h_n(x) - h_n(y) \xrightarrow{n \to \infty} h(x) - h(y), \quad x - y \quad \mu \times \mu \text{-a.e.}
\]
Using (3.9), we get
\[
\left| \frac{h_n(x) - h_n(y)}{x - y} \right| = \left| \frac{h(x) - h(y)}{x - y} \Theta_n(|x|) + \frac{\Theta_n(|x|) - \Theta_n(|y|)}{x - y} h(y) \right| \leq \left| \frac{h(x) - h(y)}{x - y} \right| + |h(y)|.
\]
From (3.8) we find that both terms on the right-hand side are bounded, and therefore $\mu \times \mu$-integrable. The Dominated Convergence Theorem can be applied, yielding
\[
\int \int \frac{h_n(x) - h_n(y)}{x - y} d\mu(x) d\mu(y) \xrightarrow{n \to \infty} \int \int \frac{h(x) - h(y)}{x - y} d\mu(x) d\mu(y),
\]
and the equation (3.7) for $h$ follows from this last limit and (3.10), keeping in mind that (3.7) is valid for each $h_n$. □

Lemma 3.5. Let $\mu$ be a $\varphi$-critical measure. If $z \in \mathbb{C}$ is such that
\[
\int \frac{d\mu(x)}{|x - z|} < +\infty,
\]
then (3.7) is valid for every function of the form
\[
h(x) = \frac{g(x)}{x - z}, \quad g \in C^2_c.
\]

Proof. The same idea used to prove Lemma 3.4 can also be used here, namely approximating $h$ by a sequence of functions for which (3.7) is valid. But now the approximating sequence needs to be modified. By translating $x \mapsto x - z$ we can assume $z = 0$, and then we rewrite
\[
h(x) = \frac{g(x) x}{|x|^2}.
\]
Define
\[
h_n(x) = \frac{g(x) x}{|x|^2 + \epsilon_n^2},
\]
where $(\epsilon_n)$ is a sequence of positive numbers converging to zero. Clearly $h_n \in C^2_c$, and so equation (3.7) is valid for it.

Note that (3.11) assures us that $h$ is $\mu$-integrable. Proceeding then similarly as in the proof of Lemma 3.4
\[
\int V'(x) h_n(x) d\mu(x) \xrightarrow{n \to \infty} \int V'(x) h(x) d\mu(x),
\]
and also
\[
\frac{h_n(x) - h_n(y)}{x - y} \xrightarrow{n \to \infty} \frac{h(x) - h(y)}{x - y}, \quad \mu \times \mu\text{-a.e.}
\]
Moreover,
\[
|h_n(x) - h_n(y)| \leq \frac{|y|^2 |\mathfrak{g}(x) - |x|^2 \mathfrak{g}(y)|}{(|x|^2 + \epsilon_n^2)(|y|^2 + \epsilon_n^2)} + \frac{\epsilon_n^2 |\mathfrak{g}(x) - \mathfrak{g}(y)|}{(|x|^2 + \epsilon_n^2)(|y|^2 + \epsilon_n^2)}
\]
\[
\leq \frac{|y|^2 |\mathfrak{g}(x) - |x|^2 \mathfrak{g}(y)|}{|x|^2 |y|^2}
\]
\[
+ \left| \frac{\mathfrak{g}(x) - \mathfrak{g}(y)}{|x||y|} \right| \frac{1}{\left( \frac{|x|}{\epsilon_n} + \frac{\epsilon_n}{|x|} \right) \left( \frac{|y|}{\epsilon_n} + \frac{\epsilon_n}{|y|} \right)}
\]
\[
\leq \frac{y g(x) - x g(y)}{xy} + \left| \frac{\mathfrak{g}(x) - \mathfrak{g}(y)}{xy} \right|,
\]
where in the last inequality we used $t + 1/t > 1$ for $t \geq 0$.\]
Now, using the trivial decompositions
\[
yg(x) - xg(y) = \frac{1}{xy}(g(x) - g(y)) - \frac{g(x)}{xy},
\]
\[
xy(x - y) - yg(x) - yg(y) = \frac{1}{xy}(g(x) - g(y)) + \frac{g(x)}{xy}
\]
in (3.14), we get
\[
\left| \frac{h_n(x) - h_n(y)}{x - y} \right| \leq \left( \frac{1}{|x|} + \frac{1}{|y|} \right) \left| \frac{g(x) - g(y)}{x - y} \right| + \frac{2}{|y|} \left| \frac{g(x)}{x} \right|
\]

Since \( g \in C^2_c \), the quotient
\[
\frac{g(x) - g(y)}{x - y},
\]
is bounded, as well as \( g \), say both by \( M \). This last inequality then implies
\[
\left| \frac{h_n(x) - h_n(y)}{x - y} \right| \leq \frac{M}{|x|} + \frac{M}{|y|} + \frac{2M}{|y|}.
\]
From assumption (3.11) the right-hand side above is \( \mu \times \mu \)-integrable. From the Dominated Convergence Theorem and (3.13) we conclude
\[
\int \int \frac{h_n(x) - h_n(y)}{x - y} d\mu(x) d\mu(y) \stackrel{n \to \infty}{\to} \int \int \frac{h(x) - h(y)}{x - y} d\mu(x) d\mu(y).
\]

Since (3.7) is valid for the \( h_n \)'s, this last limit together with (3.12) implies (3.7) for \( h \).

A combination of Lemmas 3.4 and 3.5 leads to the following.

**Corollary 3.6.** Let \( \mu \) be a \( \varphi \)-critical measure. Suppose \( z_1, z_2, \ldots, z_{N - 1} \) are distinct points such that
\[
\int \frac{d\mu(x)}{|x - z_j|} < +\infty, \quad j = 1, \ldots, N - 1.
\]
Then equation (3.7) is valid for the function
\[
h(x) = \prod_{j=1}^{N-1} \frac{1}{x - z_j}.
\]

**Proof.** Let \( \delta = \frac{1}{4} \min\{|z_j - z_k| : 1 \leq j < k \leq N - 1\} \). For each \( j = 1, \ldots, N - 1 \), we can take a \( C^2_c \) function \( \chi_j : \mathbb{C} \to [0, 1] \) supported in \( D(z_j, 2\delta) \) with \( \chi_j \equiv 1 \) in \( D(z_j, \delta) \). Then \( h_j = \chi_j/h \) satisfies the condition of Lemma 3.5 and \( h - \sum_{j=1}^{N-1} h_j \) satisfies the condition of Lemma 3.4. Thus the equality (3.7) holds for these functions, and then by linearity it also holds for \( h \). \( \square \)

Given a finite positive Borel measure \( \mu \), consider the function \( \mathbb{C} \to [0, +\infty] 
\[
z \mapsto \int \frac{d\mu(x)}{|x - z|}.
\]
It is the convolution of a finite Borel measure with the function \( x \mapsto \frac{1}{|x|} \), which is in \( L^1_{\text{loc}}(\mathbb{C}, m_2) \), where \( m_2 \) is the planar Lebesgue measure. Thus Tonelli’s Theorem
tells us it also belongs to $L_{loc}^1(\mathbb{C}, m_2)$, being then finite $m_2$-a.e. This implies that the Cauchy Transform $C^\mu$ of $\mu$,

\[ C^\mu(z) = \int \frac{d\mu(x)}{x - z}, \quad z \in \mathbb{C}, \]

is well-defined and finite $m_2$-a.e., and

\[ C^\mu(z) = O(z^{-1}), \quad z \to \infty. \tag{3.16} \]

As observed by Martínez-Finkelshtein and Rakhmanov [29, Lemma 5.1], an essential feature of critical measures is the fact that their Cauchy Transform satisfies an algebraic equation of degree 2.

**Proposition 3.7.** Let $\mu$ be a $\varphi$-critical measure. Then there exists a polynomial $R$ of degree $2N - 2$ such that

\[ \left( C^\mu(z) + \frac{1}{2} V'(z) \right)^2 = R(z), \quad m_2\text{-a.e.} \tag{3.17} \]

where $m_2$ is the Lebesgue measure on $\mathbb{C}$.

As we will see later, one of the consequences of Proposition 3.7 is that $\varphi$-critical measures are always compactly supported, see Proposition 3.9 below. However, since we cannot assume this a priori, our proof, at the technical level, is different from the one given in [29, Proof of Lemma 5.1], although the key ideas are similar.

**Proof.** Since $\int \frac{d\mu(x)}{|x - z|}$ is finite $m_2$-a.e., we can fix $N - 2$ distinct points $z_1, \ldots, z_{N-2}$ for which

\[ \int \frac{d\mu(x)}{|x - z_j|} < \infty, \quad j = 1, \ldots, N - 2. \tag{3.18} \]

Take $z \in \mathbb{C} \setminus \{z_1, \ldots, z_{N-2}\}$. Since we want to have \[(3.17)\] $m_2$-a.e. and $\int \frac{d\mu(x)}{|x - z|}$ is finite $m_2$-a.e., it is enough to prove the identity \[(3.17)\] under the assumption

\[ \int \frac{d\mu(x)}{|x - z|} < \infty. \tag{3.19} \]

Define

\[ h(x) = \frac{A(x)}{x - z}, \quad A(x) = \prod_{j=1}^{N-2} (x - z_j)^{-1}. \tag{3.20} \]

Then $h$ satisfies the conditions of Corollary 3.6 (with $z = z_{N-1}$) and so \[(3.17)\] is valid for this $h$. We can write

\[ \int V'(x)h(x) d\mu(x) = A(z)V'(z)C^\mu(z) + D_1(z), \tag{3.21} \]

where $D_1$ is the rational function

\[ D_1(z) = \int \frac{A(x)V'(x) - A(z)V'(z)}{(x - z)} d\mu(x), \]

whose only possible poles are the points $z_1, \ldots, z_{N-2}$, all of them simple, and

\[ D_1(z) = O(1), \quad z \to \infty. \tag{3.22} \]

On the other hand, we can write

\[ \frac{h(x) - h(y)}{x - y} = \frac{(x - y)A(z) + (z - x)A(y) + (y - z)A(x)}{(x - y)(x - z)(y - z)} - \frac{A(z)}{(x - z)(y - z)}. \]
The first term in the right-hand side is bounded away from $z_1, \ldots, z_{N-2}$, and since we are assuming (3.18) this term is $d\mu(x) \times d\mu(y)$-integrable. The second term is also $d\mu(x) \times d\mu(y)$-integrable because of assumptions (3.18) and (3.19). From this we conclude

$$\int \int \frac{h(x) - h(y)}{x - y} d\mu(x) d\mu(y) = -A(z)(C^\mu(z))^2 + D_2(z),$$

where

$$D_2(z) = \int \int \frac{(x - y)A(z) + (z - x)A(y) + (y - z)A(x)}{(x - y)(x - z)(y - z)} d\mu(x) d\mu(y),$$

is also a rational function with simple poles at $z_1, \ldots, z_{N-2}$ and no other poles. Moreover,

$$D_2(z) = O(z^{-1}), \quad z \to \infty. \tag{3.24}$$

As (3.7) is valid for this $h$, equations (3.21) and (3.23) give us

$$-A(z)(C^\mu(z))^2 + D_2(z) = A(z)V'(z)C^\mu(z) + D_1(z),$$

which is equivalent to (3.17) if we set

$$R(z) = \frac{D_2(z) - D_1(z)}{A(z)} + \frac{1}{4} (V'(z))^2. \tag{3.25}$$

Note that the poles of $A$ cancel out the possible poles of $D_1 - D_2$, and due also to the behavior of $D_1$ and $D_2$ at infinity given in (3.22), (3.24) we see that $R$ is indeed a polynomial of degree $2N - 2$. \qed

From the polynomial $R$ in (3.17), we can recover $\mu$.

**Proposition 3.8.** Suppose $\mu$ is a measure on $\mathbb{C}$ for which there exist polynomials $Q$ and $R$ such that

$$(C^\mu(z) + Q(z))^2 = R(z), \quad m_2\text{-a.e.} \tag{3.26}$$

Then $\mu$ is supported on a union of analytic arcs, which are maximal trajectories of the quadratic differential $-R(z)\,dz^2$. Moreover, in the interior of any arc of supp $\mu$, the measure $\mu$ is absolutely continuous with respect to the arclength measure, with density given by

$$d\mu(s) = \frac{1}{\pi i} \sqrt{R(s)} \, ds, \tag{3.27}$$

where $ds$ is the complex line element, chosen according to a fixed orientation of the arcs of supp $\mu$.

Observe that by Proposition 3.7, the condition (3.26) is valid for $\varphi$-critical measures with the choice $Q(z) = \frac{V'(z)}{2}$.

Since Proposition 3.8 is essentially about local properties of the measure $\mu$, it is not important if the measure $\mu$ has unbounded support or not. So the proof presented here is essentially the same as the one given in [29, Lemma 5.2] for measures with bounded support, which in turn is modelled after a result in [8]. We decided to give a full detailed proof here for the benefit of the reader, and also to correct a sign misprint in [29]: in the first centered formula following formula (5.23) in [29] the right-hand side should have a $-$-sign.
Proof. Let $G$ be a simply connected domain not containing zeros of $R$ and intersecting $\text{supp} \mu$. Fix a point $z_0 \in \text{supp} \mu \cap G$ and select a branch of $\sqrt{R}$ in $G$, and define
\[
\xi(z) = \int_{z_0}^{z} \sqrt{R(s)} \, ds, \quad z \in G.
\]
Reducing $G$ if necessary, we can assume $\xi$ is a conformal mapping between $G$ and
\[
\xi(G) = \mathbb{I} \times i \mathbb{J} = \hat{G}, \quad \mathbb{I}, \mathbb{J} \subset \mathbb{R}
\]
intervals both containing $z = 0$.

For this same branch of $\sqrt{R}$, define for $m_2$-a.e. $z \in G$ a function $\chi$ by the formula
\[
\chi(z) = C^\mu(z) + Q(z) \sqrt{R(z)}.
\]
Due to (3.26) $\chi$ assumes values in $\{-1, 1\}$. A simple application of the Cauchy-Green Formula \[22, \text{pg. 491}\] gives us
\[
\frac{\partial C^\mu}{\partial \xi} = -\pi d\mu
\]
in the sense of distributions, and we conclude (also in the sense of distributions)
\[
\frac{\partial}{\partial \xi} \left( \sqrt{R(\xi)} \chi(\xi) \right) = \sqrt{R(\xi)} \frac{\partial \chi}{\partial \xi}(\xi) = -\pi d\mu(\xi). \quad (3.28)
\]

Now, $\xi$ is a conformal map from $G$ onto $\hat{G}$, with inverse denoted by $z = z(\xi)$, which induces a map between distributions in the $z$-variable and the $\xi$-variable, say $u \mapsto u_*$, where $u_*$ is the distribution acting on test functions in $\hat{G}$ via the formula
\[
\langle u, \phi \rangle = \langle u_*, \psi \rangle,
\]
where $\psi(\xi) = \phi(z)$.

Now we will calculate the (distributional) derivative of $\chi_*(\xi) = \chi(z(\xi))$. If $\mu_*$ is the pushforward measure of $\mu$ induced by $\xi$ in $\hat{G}$ then (3.28) and the chain rule applied to $\psi(\xi) = \phi(z)$ imply
\[
\langle \mu_*, \psi \rangle = \langle \mu, \phi \rangle.
\]

Using the Change of Variables Formula for the Lebesgue measure
\[
\langle \mu_*, \psi \rangle = \frac{1}{\pi} \int \left| \frac{\partial \psi}{\partial \xi}(\xi) \right| \chi_*(\xi) \left| z'(\xi) \right|^2 \, dm_2(\xi),
\]

since $\frac{\partial z}{\partial \xi} = 0$ and $\frac{\partial \xi}{\partial z} = \sqrt{R(z)}$. Thus by the Change of Variables Formula for the Lebesgue measure
\[
\langle \mu_*, \psi \rangle = \frac{1}{\pi} \int \left| \frac{\partial \psi}{\partial \xi}(\xi) \right| \left( \sqrt{R(\xi)} \right) \chi_*(\xi) \left| z'(\xi) \right|^2 \, dm_2(\xi)
\]

Due to (3.26) $\chi$ assumes values in $\{-1, 1\}$. A simple application of the Cauchy-Green Formula \[22, \text{pg. 491}\] gives us
\[
\frac{\partial C^\mu}{\partial \xi} = -\pi d\mu
\]
in the sense of distributions, and we conclude (also in the sense of distributions)
\[
\frac{\partial}{\partial \xi} \left( \sqrt{R(\xi)} \chi(\xi) \right) = \sqrt{R(\xi)} \frac{\partial \chi}{\partial \xi}(\xi) = -\pi d\mu(\xi). \quad (3.28)
\]

Now, $\xi$ is a conformal map from $G$ onto $\hat{G}$, with inverse denoted by $z = z(\xi)$, which induces a map between distributions in the $z$-variable and the $\xi$-variable, say $u \mapsto u_*$, where $u_*$ is the distribution acting on test functions in $\hat{G}$ via the formula
\[
\langle u, \phi \rangle = \langle u_*, \psi \rangle,
\]
where $\psi(\xi) = \phi(z)$.

Now we will calculate the (distributional) derivative of $\chi_*(\xi) = \chi(z(\xi))$. If $\mu_*$ is the pushforward measure of $\mu$ induced by $\xi$ in $\hat{G}$ then (3.28) and the chain rule applied to $\psi(\xi) = \phi(z)$ imply
\[
\langle \mu_*, \psi \rangle = \langle \mu, \phi \rangle.
\]

Using the Change of Variables Formula for the Lebesgue measure
\[
\langle \mu_*, \psi \rangle = \frac{1}{\pi} \int \left| \frac{\partial \psi}{\partial \xi}(\xi) \right| \left( \sqrt{R(\xi)} \right) \chi_*(\xi) \left| z'(\xi) \right|^2 \, dm_2(\xi)
\]

Since $\frac{\partial z}{\partial \xi} = 0$ and $\frac{\partial \xi}{\partial z} = \sqrt{R(z)}$. Thus by the Change of Variables Formula for the Lebesgue measure
\[
\langle \mu_*, \psi \rangle = \frac{1}{\pi} \int \left| \frac{\partial \psi}{\partial \xi}(\xi) \right| \left( \sqrt{R(\xi)} \right) \chi_*(\xi) \left| z'(\xi) \right|^2 \, dm_2(\xi)
\]

Since $\frac{\partial z}{\partial \xi} = 0$ and $\frac{\partial \xi}{\partial z} = \sqrt{R(z)}$. Thus by the Change of Variables Formula for the Lebesgue measure
\[
\langle \mu_*, \psi \rangle = \frac{1}{\pi} \int \left| \frac{\partial \psi}{\partial \xi}(\xi) \right| \left( \sqrt{R(\xi)} \right) \chi_*(\xi) \left| z'(\xi) \right|^2 \, dm_2(\xi)
\]
where we used that $z'(ξ) = \frac{1}{ξ'(z)} = \frac{1}{\sqrt{R(z)}}$. The result is that $\langle μ_*, ψ \rangle = -\frac{1}{π} \left\langle \frac{∂χ_*}{∂ξ}, ψ \right\rangle$ which can be rewritten as

$$\frac{∂χ_*}{∂ξ} = \frac{1}{2} \frac{∂χ_*}{∂x} + i \frac{∂χ_*}{∂y} = -πdμ_*,$$

(3.29)

where $ξ = x + iy$ with $x, y ∈ \mathbb{R}$. As $μ$ is a real measure, we conclude $\frac{∂χ_*}{∂y} = 0$, which means $χ_*(ξ) = g(\text{Re} ξ)$, for a real function $g$ defined on the interval $I$.

If $λ$ is the pushforward measure of $μ_*$ induced by $ξ → \text{Re} ξ$, using (3.29) we can conclude in a similar fashion that

$$\frac{dg}{dx} = -αdλ,$$

where $α$ is a positive constant, and so

$$g(x) = β - α \int_{-∞}^{x} dλ,$$

for some real constant $β$. But $g$ only assumes values in $\{1, -1\}$, and since $\int_{-∞}^{x} dλ$ is non-decreasing, we conclude $g$ is non-increasing and also that there exists $x_0$ in $I$ for which $λ = cδ_{x_0}$, where $c$ is a positive constant, and then

$$g(x) = 1 - 2 \int_{-∞}^{x} dδ_{x_0}.$$

In terms of $μ_*$ and $χ_*$, this means that there exists a vertical segment $L = \{x_0\} × iJ$ in $\hat{G}$ with $χ_*$ equal to 1 at the left side of $L$ and $-1$ at the right side of $L$, and supp $μ_*$ is contained in $L$. But for such $χ_*$ a direct computation shows

$$\frac{∂χ_*}{∂ξ} = \frac{1}{2} \frac{∂χ_*}{∂x} = -dy,$$

where $dy$ is the Lebesgue measure on $L$, so by (3.29)

$$dμ_* = \frac{1}{πi} dy = \frac{1}{iπ} dt.$$

where $dt$ is the line element in $L$, oriented such that $χ_* = 1$ on the positive side of $L$. Since $ξ$ is a conformal map we can pullback the measures $dμ_*$ and $dt$, obtaining

$$dμ(s) = \frac{1}{πi} \sqrt{R(s)} ds,$$

where $ds$ is the complex line element on the trajectory $γ = ξ^{-1}(L)$ chosen with orientation such that $ξ(z) = 1$ on the positive side (i.e., on the left-hand side) of $γ$.

Finally, the trajectory arcs of supp $μ$ must be maximal, because the arguments above show that supp $μ$ is an analytic arc in a neighborhood of any point of its support that is not a zero of $R$. □

For the next proposition, it is important that $φ = \text{Re} V$ with $V$ polynomial.

**Proposition 3.9.** For any $φ$-critical measure $μ$, the support supp $μ$ is compact and can be written as

$$\text{supp } μ = \bigcup_j α_j,$$

(3.30)

where each $α_j$ is an analytic arc and the union is taken over a finite set of indices. Each $α_j$ is a critical trajectory of the quadratic differential $-R(z)dz^2$. 
Proof. Propositions 3.7 and 3.8 tell us that the measure $\mu$ is supported on maximal trajectories of the quadratic differential

$$-R(z)dz^2 = -\left(C^\mu(z) + \frac{1}{2} V'(z)\right)^2 dz^2,$$

which has a pole of order $2N + 2$ at $z = \infty$. The general theory of quadratic differentials assures us that there exists a neighborhood $U$ of $z = \infty$ such that every trajectory intersecting $U$ ends up at $z = \infty$ in at least one direction [44, Theorem 7.4]. In particular, if $\gamma \cap U \neq \emptyset$ for some maximal trajectory $\gamma$ of $\text{supp}\mu$, then $\gamma$ would be unbounded, and due to (3.27) the measure $\mu$ would not be finite.

Since $R$ is a polynomial, the only bounded trajectories are critical trajectories that connect two different zeros of $R$. Since there are only finitely many zeros of $R$, the union in (3.30) is certainly over a finite number of arcs. □

It is worth noting that the expression (3.25) for $R$ in terms of the function $A$ defined in (3.20), depends on the chosen points $z_1, \ldots, z_{N-2}$ at which the Cauchy transform of $\mu$ is convergent. Since we now know by Proposition 3.9 that $\text{supp}\mu$ is compact, we can take the constant function $A(x) = 1$ in (3.20), and the rest of the proof of (3.17) then works fine for this $A$. Repeating all the computation, we also end up with a nicer representation for $R$,

$$R(z) = \left(\frac{V'(z)}{2}\right)^2 - \int \frac{V'(x) - V'(z)}{x - z} d\mu(x), \quad (3.31)$$

which is well-known for equilibrium problems in polynomial external fields on the real line, see e.g. [17]. This representation also allows us to derive some extra properties for $R$. For example, from (3.31) one easily sees that $R(z) - \frac{(V'(z))^2}{4}$ is a polynomial of degree $n - 2$, whose leading coefficient coincides with minus the leading coefficient of $V'$.

Also, if $V$ is quadratic, say

$$V(z) = az^2 + bz + c, \quad a \neq 0,$$

then (3.31) reduces to

$$R(z) = \left(\frac{2az + b}{2}\right)^2 - 2a,$$

since $\mu$ is a probability measure. Then the zeros of $R$ are

$$z_\pm = \frac{-b \pm \sqrt{2a}}{2a},$$

for some choice of the branch of the square root, and the support of the $\varphi$-critical measure is the segment from $z_-$ to $z_+$. For $V$ cubic, the polynomial $R(z) - \frac{(V'(z))^2}{4}$ is linear, so there is one free parameter to be determined by the requirement of extra conditions, see for instance [11,16,24]. See also [15] for a similar situation.

3.3. Critical measures and the S-property. By Proposition 3.8 the support of a $\varphi$-critical measure $\mu$ consists of analytic arcs and their endpoints. The analytic arcs are trajectories of the quadratic differential $-R(z)dz^2$. Each regular point $z \in \text{supp}\mu$ has a neighborhood, say $G$, for which $G \cap \text{supp}\mu$ is an analytic arc. The only non-regular points are the zeros of $R$.

The next result is taken from [29].
Proposition 3.10. Let $\mu$ be a $\varphi$-critical measure, with $\text{supp} \mu = \bigcup \alpha_j$ as in Proposition 3.7.

i) Then there exist constants $\omega_j \in \mathbb{R}$ such that

$$U^\mu(z) + \frac{1}{2} \varphi(z) = \omega_j, \quad z \in \alpha_j.$$ 

ii) If $z \in \text{supp} \mu$ is regular, then

$$\frac{\partial}{\partial n_+} \left( U^\mu + \frac{1}{2} \varphi \right)(z) = \frac{\partial}{\partial n_-} \left( U^\mu + \frac{1}{2} \varphi \right)(z),$$

where $n_\pm$ are the normal vectors to $\text{supp} \mu$ at $z$ pointing in opposite directions.

Proof. See Lemma 5.4 of [29]. □

An immediate consequence of part ii) of Proposition 3.10 is.

Corollary 3.11. If $V$ is a polynomial, $\varphi = \text{Re} V$ and $\Gamma \subset \mathbb{C}$ is a contour satisfying the growth condition (2.3) and whose equilibrium measure $\mu^{\varphi, \Gamma}$ is a $\varphi$-critical measure, then $\Gamma$ has the S-property in the external field $\varphi$.

As a final remark, we notice that the cornerstone of the present section is the Proposition 3.7 which gives us that the Cauchy transform of a $\varphi$-critical measure is an algebraic function. One could also analyse (3.17) from another perspective, namely given $V'$, find a polynomial $R$ and a probability measure $\mu$ such that the algebraic equation (3.17) is satisfied. This is the spirit of some works already present in the literature, e.g. [23, 29, 37, 38]. It is expected that the pairs $(R, \mu)$ satisfying (3.17) depend on continuous parameters, but as it follows from the present work, only for a finite number of pairs $(R, \mu)$ the measure $\mu$ is an equilibrium measure (of some admissible contour) in the external field $\varphi = \text{Re} V$. More precisely, just for a finite number of pairs $(R, \mu)$ obtained in this way, all the constants $\omega_j$ appearing in Proposition 3.10 i) coincide. In such a case there is a choice of partition $P$ and a contour $T(P)$ for which $\mu$ is the equilibrium measure of $\Gamma$ in the external field $\varphi$.

4. Critical sets

We continue to use the notation introduced in the beginning of Section 3. So for $h \in C^2$ we have $h_t$ as in (3.2), and we denote by $\mu^t$ the pushforward of a measure $\mu$ induced by $h_t$. For a subset $F \subset \mathbb{C}$, we denote $F^t = h_t(F)$.

In Section 3 we studied variations of the energy functional $I^\varphi$ induced by $h$, viewed as acting on measures, see (3.3). But we might as well study variations of $I^\varphi$ when viewed as acting on sets, that is, we might consider the limit

$$D_h I^\varphi(F) = \lim_{t \to 0} \frac{I^\varphi(F^t) - I^\varphi(F)}{t}.$$ 

(4.1)

Just as the vanishing of $D_h I^\varphi(\mu)$ defines critical measures, the vanishing of $D_h I^\varphi(F)$ defines critical sets.

Definition 4.1. We say that a closed set $F$ is a $\varphi$-critical set (or simply a critical set) if the limit in (4.1) exists for every $h \in C^2$, and

$$D_h I^\varphi(F) = 0$$

for every $h \in C^2$. 
The remarkable fact is that the limit (4.1) coincides with the derivative of $I^\varphi$ at the equilibrium measure in external field of $F$.

**Proposition 4.2.** Let $F \subset \mathbb{C}$ be a closed set with $-\infty < I^\varphi(F) < +\infty$. Then the limit in (4.1) exists and

$$D_h I^\varphi(F) = D_h I^\varphi(\mu^{\varphi,F}).$$

(4.2)

The existence of the equilibrium measure in the external field $\mu^{\varphi,F}$ is guaranteed by Theorem 2.5. Proposition 4.2 is due to Rakhmanov, who gives a proof in [33, Section 9.10] and refers to an unpublished manuscript where the argument is first given. We found that one of the steps in the proof in [33], namely the equivalent of (4.11) below, might require some additional explanation, and therefore we decided to present here a detailed proof.

To prove Proposition 4.2 we will need an auxiliary lemma. We use $\mathcal{M}_0$ to denote the set of signed measures defined on $\mathbb{C}$, whose positive and negative parts belong to $\mathcal{M}_1(\mathbb{C})$ (see the beginning of Section 2.1 for the definition of $\mathcal{M}_1(\mathbb{C})$). Recall that a sequence of finite signed measures $(\sigma_n)$ converges vaguely to a finite signed measure $\sigma$ if

$$\int f d\sigma_n \xrightarrow{n \to \infty} \int f d\sigma,$$

for every continuous functions $f$ with compact support.

**Lemma 4.3.** The functional

$$\mathcal{M}_0 \ni \sigma \mapsto I(\sigma) = \iint \log \frac{1}{|x-y|} d\sigma(x) d\sigma(y)$$

is well-defined on $\mathcal{M}_0$ and strictly positive for $\sigma \neq 0$. If $(\sigma_n)$ is a sequence of signed measures in $\mathcal{M}_0$ with

$$I(\sigma_n) \xrightarrow{n \to \infty} 0,$$

then the sequence $(\sigma_n)$ converges vaguely to the null measure.

**Proof.** The lemma is well known if we consider signed measures supported in the unit disc, see for example [27, pg. 80, Theorem 1.16, pg. 88, Lemma 1.3]. The main issue here is that we are dealing with signed measures with possibly unbounded support. The proof we are going to give is actually obtained by collecting known results together and fitting them to our needs here.

Given two measures $\mu, \nu$ with finite energies $I(\mu), I(\nu)$, the integral

$$I(\mu, \nu) = \iint \log \frac{1}{|x-y|} d\mu(x) d\nu(y),$$

is well defined and finite [31, Theorem 2.2], that is, $|\log |x-y|^{-1}| \in L^1(\mu \times \nu)$. In particular, if $\sigma = \sigma_+ - \sigma_- \in \mathcal{M}_0$, then the value

$$I(\sigma) = I(\sigma_+) + I(\sigma_-) - 2I(\sigma_+, \sigma_-)$$

is also well defined and finite. Due to a result of Mattner [31, Theorem 2.2] it is positive, i.e., $I(\sigma) \geq 0$ and $I(\sigma) = 0$ if and only if $\sigma$ is the zero measure.

For $\sigma \in \mathcal{M}_0$ compactly supported and absolutely continuous with respect to the Lebesgue measure $m_2$, say $d\sigma = f dm_2$, with $f \in C_\infty$ and real-valued satisfying $\int f dm_2 = 0$, Cegrell et al. [13, Lemma 2.4] obtained the representation

$$I(\sigma) = 2\pi \int \frac{|\hat{f}(x)|^2}{|x|^2} dm_2(x),$$

(4.3)
for $I(\sigma)$, where

$$\hat{f}(z) = \int f(x)e^{-2\pi i \Re(z \overline{x})} \, dm_2(x)$$

is the Fourier transform of $f$ (induced by $\mathbb{R}^2$, so that $\Re(z \overline{w}) = \Re (x \overline{y}) + \Im (x \overline{y})$ is just the inner product in $\mathbb{R}^2$). For $\sigma = \sigma_+ - \sigma_- \in \mathcal{M}_0$ and a given sequence of approximation to the identity $(\psi_n)$ ($\psi_n$ is smooth and independent of $\sigma$), satisfying

$$\int \psi_n \, dm_2 = 1, \quad \psi_n(z) \geq 0, \quad \text{for} \ z \in \mathbb{C}, \quad \text{and} \quad \psi_n(z) \equiv 0 \quad \text{for} \ |z| > \frac{1}{n},$$

Cegrell et al. also constructed a sequence of absolutely continuous signed measures $\mu_n = \psi_n \ast \lambda_n \in \mathcal{M}_0$ with $\lim_{n \to \infty} I(\mu_n) = I(\sigma)$, where

$$\lambda_n = \sigma |_{\Sigma_n} - \sigma(\Sigma_n) \lambda,$$ \hspace{1cm} (4.4)

and $\lambda$ is the normalized Lebesgue measure on the unit disc and $(\Sigma_n)$ a sequence of discs $\Sigma_n = D(0, R_n)$ whose respective sequence of radii $(R_n)$ converges to $+\infty$ when $n \to \infty$.

So now consider a sequence $(\sigma_m)$ in $\mathcal{M}_0$ with $I(\sigma_m) \to 0$. For each $m$ consider the measures $\lambda_{n,m}$ defined as in (4.4) for $\sigma = \sigma_m$. For each $m$, we can take $n = n(m)$ large enough, such that for $\mu_m := \lambda_{n,m} \ast \psi_n$,

$$|I(\mu_m) - I(\sigma_m)| < \frac{1}{m}, \quad (4.5)$$

and we can make sure that $n(m+1) > n(m)$, for all $m$. The explicit form of $\mu_m$ and this last condition on the indices imply in particular

$$\mu_m - \sigma_m \to 0 \text{ vaguely.} \quad (4.6)$$

The measures $\mu_m$ are absolutely continuous w.r.t. $m_2$ with a smooth and compactly supported density, say $d\mu_m = f_m \, dm_2$. Equations (4.3) and (4.5) and the convergence $I(\sigma_m) \to 0$ then imply

$$\frac{\hat{f}_m(x)}{x} \to 0 \text{ in } L^2(m_2) \quad \text{as } m \to \infty.$$

If $h = h_1 + \i h_2 \in C^\infty$, then $h_j, f_m \in L^1(m_2) \cap L^2(m_2)$ and also $\hat{h}_j \in L^2(m_2)$, and the Plancherel Theorem and Cauchy-Schwarz inequality imply for $j = 1, 2$,

$$\left| \int h_j \, d\mu_m \right| = \left| \int h_j(x) f_m(x) \, dm_2(x) \right|$$

$$= \left| \int \hat{h}_j(z) \hat{f}_m(z) \, dm_2(z) \right|$$

$$\leq \|x \hat{h}_j\|_2 \left\| \frac{\hat{f}_m}{x} \right\|_2 \to 0, \quad \text{as } m \to \infty,$$

so that

$$\lim_{m \to \infty} \int h \, d\mu_m = 0. \quad (4.7)$$

Finally, since $C^\infty_c$ is dense in $C_c$ w.r.t. the uniform norm and $| \int f_m \, dm_2 |$ is uniformly bounded, the inequality

$$\left| \int h \, d\mu \right| \leq \| h - h_n \|_\infty \left| \int f_m \, dm_2 \right| + \left| \int h_n \, d\mu_m \right|$$
for a sequence \((h_m)\) of functions in \(C^\infty\) converging to \(h \in C_c\) in the uniform norm shows the limit (4.10) is also true for \(h \in C_c\), that is, \(\mu_m \to 0\) vaguely. In virtue of (4.6) this is equivalent to the vague convergence \(\sigma_m \to 0\).

Proof of Proposition 4.2. For \(t\) small enough and \(h_t\) as in (3.2), the inverse map \(h_t^{-1}\) is a well defined \(C^2\) function which satisfies

\[
\frac{1}{x - y} h_t^{-1}(x) - \frac{1}{x - y} h_t^{-1}(y) = 1 - t \frac{h(x) - h(y)}{x - y} + o(t),
\]

with the \(o\) terms uniform in \(x, y \in \mathbb{C}\).

For simplicity, denote by \(\mu_F\) and \(\mu_{F'}\) the equilibrium measures in the external field \(\varphi\) of \(F\) and \(F'\), respectively. By Rakhmanov’s Theorem 2.5 and the assumption \(h \in C_c\) these measures certainly exist.

Clearly, in view of the definitions (3.3) and (4.1) and the Proposition 3.1 it is enough to prove that

\[
I^\varphi(F^t) - I^\varphi(F) = I^\varphi(\mu_F^t) - I^\varphi(\mu_F) + o(t), \quad t \to 0,
\]

which comes down to proving

\[
I^\varphi(\mu_{F'}) - I^\varphi(\mu_F^t) = o(t), \quad t \to 0.
\]

For a given \(t\), denote by \(\mu_F^t\) the pushforward measure of \(\mu_{F'}\) induced by the inverse mapping \(h_t^{-1}\). Bearing in mind equations (4.8), (4.9) and mimicking the proof of Proposition 3.1 we obtain

\[
I^\varphi(\mu_F^t) - I^\varphi(\mu_{F'}) = -t D_h I^\varphi(\mu_{F'}) + o(t), \quad t \to 0,
\]

and also from Proposition 3.1

\[
I^\varphi(\mu_F^t) - I^\varphi(\mu_F) = t D_h I^\varphi(\mu_F) + o(t), \quad t \to 0.
\]

This in turn implies

\[
0 \leq I^\varphi(\mu_F^t) - I^\varphi(\mu_{F'}) = I^\varphi(\mu_F) - I^\varphi(\mu_{F'}^t) + t(D_h I^\varphi(\mu_F) - D_h I^\varphi(\mu_{F'})) + o(t)
\]

\[
\leq t(D_h I^\varphi(\mu_F) - D_h I^\varphi(\mu_{F'})) + o(t), \quad t \to 0,
\]

so we are done if we can show

\[
\lim_{t \to 0} D_h I^\varphi(\mu_{F'}) = D_h I^\varphi(\mu_F).
\]

To obtain (4.11), assume for a moment

\[
\int f \, d\mu_{F'} \xrightarrow{t \to 0} \int f \, d\mu_F, \quad f \in C_c,
\]

that is, \(\mu_{F'} \to \mu_F\) vaguely, to be proved later. Since \(h \in C_c^2\) the function

\[
(x, y) \mapsto \frac{h(x) - h(y)}{x - y}
\]

is continuous with compact support, so that by (4.12)

\[
\int \int \frac{h(x) - h(y)}{x - y} \, d\mu_{F'}(x) \, d\mu_{F'}(y) \xrightarrow{t \to 0} \int \int \frac{h(x) - h(y)}{x - y} \, d\mu_F(x) \, d\mu_F(y).
\]
The function $hV'$ belongs also to $C^2_c$, so again by (4.12)
\[
\int h(x)V'(x) \, d\mu_F(x) \xrightarrow{t \to 0} \int h(x)V'(x) \, d\mu_F(x).
\] (4.14)

Using Proposition 3.1 and (4.13), (4.14), the limit (4.11) follows and the proof is done, provided that we have (4.12).

The proof of (4.12) will be done in two steps, namely $\mu_Ft \to \mu_F$ and $\mu_Ft - \mu_F \to 0$ vaguely.

The first of these limits follows directly from the Change of Variables Formula,
\[
\int f \, d\mu_Ft = \int f \circ h_t \, d\mu \quad t \to 0,
\]
because $f$ has compact support, and so it is bounded, and $h_t(z) \to z$ as $t \to 0$, and the Dominated Convergence Theorem can be applied.

The difference $D h I^\varphi (\mu_F) - D h I^\varphi (\mu_Ft)$ is bounded, thanks to Proposition 3.1, so (4.10) implies $I^\varphi (\mu_Ft) - I^\varphi (\mu_F) \to 0$. Now, following notations of the previous Lemma,
\[
0 \leq I(\mu_Ft - \mu_F) = 2I^\varphi (\mu_Ft) + 2I^\varphi (\mu_F) - 4I^\varphi \left( \frac{\mu_Ft + \mu_F}{2} \right)
\]
\[
\leq 2I^\varphi (\mu_Ft) + 2I^\varphi (\mu_F) - 4I^\varphi (\mu_F) = 2I^\varphi (\mu_Ft) - 2I^\varphi (\mu_F) \to 0
\] (4.15)
as $t \to 0$. The inequality in (4.15) is obtained by noting that $\frac{1}{2}(\mu_Ft + \mu_F) \in M^\varphi (F^\circ)$, so its weighted energy is larger than $I^\varphi (\mu_Ft)$. Lemma 4.3 then implies $\mu_Ft - \mu_F \to 0$ vaguely.

An immediate consequence of Proposition 4.2 is the following.

**Corollary 4.4.** The equilibrium measure $\mu_{\varphi,F}$ in the external field $\varphi$ of a $\varphi$-critical set $F$ is a $\varphi$-critical measure.

### 5. Proof of Theorem 2.3

In this section we are going to prove Theorem 2.3. Recall we are assuming that $V$ is a polynomial of degree $N$, $\varphi = \Re V$ and $\mathcal{T} = \mathcal{T}(\mathcal{P})$ is the class of admissible contours associated to a fixed noncrossing partition $\mathcal{P}$, as given in Definition 2.2. This setup defines the max-min problem $(V, \mathcal{T})$, see (2.8).

#### 5.1. The collection $\mathcal{T}_M$.

In the first step we are going to restrict the class of contours $\mathcal{T}$. The contours in $\mathcal{T}$ stretch out to infinity in certain directions in which $\varphi \to +\infty$, as specified by the partition $\mathcal{P}$. However, it is not forbidden that some parts of $\Gamma \in \mathcal{T}$ are in regions of the plane where $\varphi$ is very negative. Then also $I^\varphi (\Gamma)$ will be very negative, and so such $\Gamma$ will be far from optimal for the max-min problem (2.8).

Here we are going to make this precise, and we show that for $M$ large enough we can restrict to a certain subclass $\mathcal{T}_M$ of $\mathcal{T}$ that we are going to define first.

For $M > 0$ large enough, the level set
\[
\varphi^{-1}(-M) = \{ z \in \mathbb{C} \mid \varphi(z) = -M \}.
\] (5.1)
consists of \(N\) disjoint analytic arcs, each of them stretching out to infinity in its both ends. This is so because \(\varphi\) is the real part of a polynomial of degree \(N\). Then, if \(M > 0\) is large enough,
\[
\Delta_M = \varphi^{-1}(-\infty, -M) \cap \{z \in \mathbb{C} \mid \text{dist}(z, \varphi^{-1}(-M)) > 8\}
\]  
(5.2)
is an open non-empty set and its boundary consists of a union of \(N\) pairwise disjoint analytic arcs. The distance in (5.2) is the usual Euclidean distance
\[
\text{dist}(z, X) = \inf_{x \in X} |z - x|.
\]
The set \(\Delta_M\) is indicated by the gray regions in Figures 2 and 4. We also assume \(M\) satisfies
\[
-M < \sup_{\Gamma \in \mathcal{T}} I^\varphi(\Gamma).
\]  
(5.3)

Then for such \(M\) we consider the subclass
\[
\mathcal{T}_M = \{\Gamma \in \mathcal{T} \mid \Gamma \cap \Delta_M = \emptyset\},
\]  
(5.4)
and its closure
\[
\mathcal{F}_M = \overline{\mathcal{T}_M}
\]  
(5.5)
with respect to Hausdorff metric on closed subsets of \(\mathbb{C}\) with the hyperbolic distance. The sets in \(\mathcal{F}_M\) are closed subsets of \(\mathbb{C}\), but they are not necessarily finite unions of contours, since this property is not preserved under taking closure in the Hausdorff metric.

However, because of Definition 2.2 ii), each \(F \in \mathcal{F}_M\) has at most \(|\mathcal{P}_0|\) components that are all unbounded in \(\mathbb{C}\), since this property is preserved by taking closure in the Hausdorff metric. Also
\[
F \cap \Delta_M = \emptyset, \quad \text{for all } F \in \mathcal{F}_M.
\]  
(5.6)

**Proposition 5.1.** We have
\[
\sup_{\Gamma \in \mathcal{T} \setminus \mathcal{T}_M} I^\varphi(\Gamma) < \sup_{\Gamma \in \mathcal{T}_M} I^\varphi(\Gamma) \leq \sup_{F \in \mathcal{F}_M} I^\varphi(F) < +\infty.
\]  
(5.7)

**Proof.** The inequality in (5.7) between the supremum over \(\mathcal{T}_M\) and the supremum over \(\mathcal{F}_M\) is trivial.

We start with the proof of the first inequality in (5.7). To prove this we take \(\Gamma \in \mathcal{T} \setminus \mathcal{T}_M\). Since all connected components of \(\Gamma\) are unbounded, see item ii) of Definition 2.2 and \(\Gamma\) intersects the set \(\Delta_M\) from (5.2), we have that \(\Gamma\) contains a connected subset \(\Gamma_0\) satisfying
\[
\Gamma_0 \subset \varphi^{-1}(-\infty, -M) \quad \text{and} \quad \text{diam} \Gamma_0 \geq 8.
\]
Then \(\text{cap}(\Gamma_0) \geq \frac{\text{diam} \Gamma_0}{4} \geq 2\), which means that \(I(\Gamma_0) \leq \log \frac{1}{2}\), see for example [35, page 138]. Let \(\omega\) be the equilibrium measure on \(\Gamma_0\) (without external field). Then
\[
I(\omega) = I(\Gamma_0) \leq \log \frac{1}{2},
\]
and
\[
I^\varphi(\Gamma_0) \leq I^\varphi(\omega) = I(\omega) + \int \varphi \, d\omega \leq \log \frac{1}{2} - M,
\]
because \(\Gamma_0\) is contained in the region where \(\varphi < -M\). Then by (5.3)
\[
I^\varphi(\Gamma) \leq I^\varphi(\Gamma_0) \leq \log \frac{1}{2} + \sup_{\Gamma' \in \mathcal{T}} I^\varphi(\Gamma').
\]
Since \(\Gamma \in \mathcal{T} \setminus \mathcal{T}_M\) is arbitrary, and \(\log \frac{1}{2} < 0\), we find the first inequality in (5.7).
It remains to prove that $\sup_{F \in F_M} I^\varphi(F)$ is finite. To this end, we construct a closed curve $\gamma$, contained in $C \setminus \Delta_M$ and intersecting all the connected components of $\partial \Delta_M$, see Figure 4 for a pictorial configuration of $\gamma$.

Recall that by item ii) of Definition 2.2 the connected components of sets in $T$ are unbounded and stretch out to infinity in at least two of the sectors $S_1, \ldots, S_N$. Assuming in addition that $\Gamma \in T_M$, that is, $\Gamma \cap \Delta_M = \emptyset$, it is easy to see that $\Gamma_A \cap \gamma \neq \emptyset$ for any connected component $\Gamma_A$ of $\Gamma$. This property is preserved by taking closures with respect to Hausdorff metric, that is, it is also valid

$$ F_A \cap \gamma \neq \emptyset \tag{5.8} $$

for any connected component $F_A$ of a set $F$ belonging to $F_M$.

Fix $R > 4$ for which $\gamma \subset D_{R-4}$. A similar reasoning as the one that led to (5.8) shows that

$$ F_A \cap \partial D_R \neq \emptyset, \tag{5.9} $$

for any connected component $F_A$ of $F \in F_M$. From (5.8) and (5.9) we see that there exists a connected compact $F_0 \subset F$ with

$$ F_0 \cap \gamma \neq \emptyset, \quad F_0 \cap \partial D_R \neq \emptyset, \quad F_0 \subset D_R. $$

This means that $\text{diam} F_0 \geq 4$, because $\gamma \subset D_{R-4}$. Then $\text{cap} F_0 \geq 1$, which implies $I(F_0) \leq 0$, and then

$$ I^\varphi(F) \leq I^\varphi(F_0) \leq I(F_0) + \sup_{z \in F_0} \varphi(z) \leq \sup_{z \in D_R} \varphi(z) < +\infty. $$

Since $R$ is independent of $F \in F_M$, the inequality $\sup_{F \in F_M} I^\varphi(F) < +\infty$ follows. □
5.2. A maximizer set $F_0$. The following is an immediate consequence of Proposition 5.1.

**Corollary 5.2.** There exists $F_0 \in \mathcal{F}_M$ such that
\[
I^\varphi(F_0) = \sup_{F \in \mathcal{F}_M} I^\varphi(F). \tag{5.10}
\]

**Proof.** The class $\mathcal{F}_M$ is closed in the Hausdorff metric. The upper semicontinuity of $I^\varphi$ given by Theorem 2.6 and Proposition 5.1 give the desired result. \qed

The next lemma assures that small variations of $F_0$ are still in $\mathcal{F}_M$.

**Lemma 5.3.** Let $F_0 \in \mathcal{F}_M$ be a maximizer set given by Corollary 5.2. For every $h \in C^2_c$, there exists $t_0 > 0$ such that $F_t \in \mathcal{F}_M$, for every $t \in (-t_0, t_0)$.

**Proof.** Since $F_0 \in \mathcal{F}_M$, we already observed that $F_0 \cap \Delta_M = \emptyset$, see (5.6). The same argument that was used to prove the first inequality in Proposition 5.1 shows that $F_0 \cap \Delta_M = \emptyset$.

Let $h \in C^2_c$ be a test function. Since $h$ is compactly supported, it follows that there is a $t_0 > 0$ such that $F_t \cap \Delta_M = \emptyset$ for all $t \in (-t_0, t_0)$.

We also take $t_0$ small enough so that $x \mapsto h_t(x) = x + th(x)$ is a $C^2$-diffeomorphism on $\mathbb{C}$ for $t \in (-t_0, t_0)$.

Let $t \in (-t_0, t_0)$, and let $(\Gamma_n)$ be a sequence in $\mathcal{T}_M$ that converges to $F_0$ in Hausdorff metric. It is then an easy fact that the convergence
\[
\Gamma_n \to F_t \quad \text{as} \quad n \to \infty \tag{5.12}
\]
in the Hausdorff metric also holds true.

Take $R > 0$ such that $\text{supp} \, h \subset D_R$. Since $h_t$ is a diffeomorphism that leaves $\mathbb{C} \setminus D_R$ invariant, we then have
\[
\Gamma_n \cap (\mathbb{C} \setminus D_R) = \Gamma_n \cap (\mathbb{C} \setminus D_R), \tag{5.13}
\]
for all $n$. Since $\Gamma_n \in \mathcal{T}_M$, it does not intersect $\Delta_M$, and so by (5.13)
\[
\Gamma_n \cap (\mathbb{C} \setminus D_R) \cap \Delta_M = \emptyset. \tag{5.14}
\]

The convergence (5.12), together with the fact that $F_t$ does not intersect the closed set $\Delta_M$, see (5.11), implies that for large enough $n$,
\[
\Gamma_n \cap D_R \cap \Delta_M = \emptyset. \tag{5.15}
\]

Then by (5.14) and (5.15) we have
\[
\Gamma_n \cap \Delta_M = \emptyset
\]
for $n$ large enough, and so $\Gamma_n \in \mathcal{T}_M$ by the definition (5.4) of the class $\mathcal{T}_M$. Because of (5.12), we then conclude that $F_t \in \mathcal{T}_M = \mathcal{F}_M$. \qed

From Lemma 5.3 it follows that the maximizer $F_0$ is a $\varphi$-critical set. Thus by Corollary 4.4 its equilibrium measure in the external field $\varphi$ is a $\varphi$-critical measure and the results of Section 3 apply to this measure. Thus we obtain the following proposition that summarizes the work we did so far.

5.3. Proposition 5.4. From Corollary 5.2, Proposition 5.3, and the results of Section 3, we obtain the following proposition that summarizes the work we did so far.

**Proposition 5.4.** The equilibrium measure in the external field $\varphi$ is a $\varphi$-critical measure.
Proposition 5.4. The equilibrium measure \( \mu_0 = \mu_{\varphi,F_0} \) of \( F_0 \) is \( \varphi \)-critical. There exists a polynomial \( R \) such that

\[
\left( C^{\mu_0}(z) + \frac{1}{2} V'(z) \right)^2 = R(z), \quad m_2 \text{-a.e.} \tag{5.16}
\]

Moreover, \( \text{supp} \, \mu_0 \) is compact and consists of a finite union of analytic arcs, which are maximal trajectories of the quadratic differential \( -R(z)dz^2 \), with each trajectory connecting two distinct zeros of \( R \). The measure \( \mu_0 \) is absolutely continuous with respect to the arclength measure on those arcs and

\[
d\mu_0(s) = \frac{1}{\pi i \sqrt{R(s)}} \, ds \tag{5.17}
\]

where \( ds \) is the complex line element. At any point \( z \in \text{supp} \, \mu_0 \) which is not a zero of \( R \), its logarithmic potential satisfies

\[
\frac{\partial}{\partial n_+} \left( U^{\mu_0} + \frac{1}{2} \varphi \right)(z) = \frac{\partial}{\partial n_-} \left( U^{\mu_0} + \frac{1}{2} \varphi \right)(z), \tag{5.18}
\]

where \( n_\pm \) are the unit normal vectors to \( \text{supp} \, \mu_0 \) at \( z \) pointing in opposite directions.

Thus by (5.18) we have that the support of \( \mu_0 \) satisfies the S-property in the external field \( \varphi \).

5.3. Proof of Theorem 2.3. Having in hands the equilibrium measure in the external field \( \mu_0 = \mu_{\varphi,F_0} \) of the maximizer set \( F_0 \) given by Corollary 5.2, our final task is to construct a contour \( \Gamma_0 \in \mathcal{T} \) for which \( \mu_{\varphi,\Gamma_0} = \mu_0 \). Note that we know that the support of \( \mu_0 \) consists of a finite union of analytic arcs, but we do not know that the full set \( F_0 \) is a contour. There is actually no reason why this should be the case, and so we will modify \( F_0 \) outside of the support of its equilibrium measure in the external field, while preserving the equilibrium measure in the external field. We also have to show that the contour \( \Gamma_0 \) that we obtain this way belongs to the class \( \mathcal{T} \).

To this end, we consider the set

\[
\Lambda = \left\{ z \in \mathbb{C} \mid U^{\mu_0}(z) + \frac{1}{2} \varphi(z) > l_0 \right\}, \tag{5.19}
\]

for the variational constant \( l_0 \) of \( \mu_0 \) appearing in the Euler-Lagrange variational conditions (2.4) and (2.5) for \( \mu_0 \).

Lemma 5.5.

For any \( \epsilon > 0 \), there exists \( R_\epsilon > 0 \) such that for every \( R \geq R_\epsilon \) the following holds. The set \( \Lambda \setminus D_R \) has exactly \( N \) connected components, each of them contained in precisely one of the sectors

\[
\left\{ z \in \mathbb{C} \mid |z| \geq R, \arg z - \theta_j \leq \frac{\pi}{2N} + \epsilon \right\}, \quad j = 1, \ldots, N, \tag{5.20}
\]

and containing in its interior the half-rays

\[
L_{j,R} = \left\{ z \in \mathbb{C} \mid |z| \geq R, \arg z = \theta_j \right\}, \quad j = 1, \ldots, N. \tag{5.21}
\]

Recall that the angles \( \theta_j \) are given in (2.7) and they are determined by the polynomial \( V \).
The boundary of $\Lambda$ consists of level curves of $U^{\mu_0} + \frac{1}{2} \varphi$, and these are exactly the trajectories of the quadratic differential $-R(z)dz^2$ given in (5.16). The quadratic differential has a pole of order $2N + 2$ at $z = \infty$. Lemma 5.5 then follows from the local behavior of trajectories of a quadratic differential near a pole of order $\geq 4$, see [44, Theorem 7.4].

Fix a number $0 < \epsilon < \frac{\pi}{2N}$ (5.22) and also $R > R_\epsilon$ for which

$$\text{supp} \mu_0 \subset D_R.$$  (5.23)

For $j = 1, \ldots, N$, let $\Lambda(j)$ be the connected component of $\overline{\Lambda} \cup \text{supp} \mu_0$ for which

$$L_{j,R} \subset \Lambda(j).$$  (5.24)

Then $\Lambda(j)$ is pathwise connected, and we can connect any point in $\Lambda(j)$ to infinity through a curve entirely contained in $\Lambda(j)$ and stretching out to infinity in the sector $S_j$ (see (2.7) for the definition of $S_j$).

The sets $\Lambda(j)$ for $j = 1, \ldots, N$ need not be disjoint. It can happen that $\Lambda(j) = \Lambda(k)$ for some pair of distinct numbers $j, k$. Of course, if their intersection is non-empty then $\Lambda(j) = \Lambda(k)$.

The next lemma is the key for completing the proof of Theorem 2.3.

**Lemma 5.6.** If $A \in P_0$ and $j, k \in A$, then $\Lambda(j) = \Lambda(k)$.

The proof of Lemma 5.6 is postponed to Section 5.4. We first show how to complete the proof of Theorem 2.3 assuming the lemma.

**Proof of Theorem 2.3.** Let $\epsilon$ and $R > R_\epsilon$ be given as in (5.22) and (5.23). For the angles $\theta_1, \ldots, \theta_N$ in (2.7), denote

$$z_j = Re^{i\theta_j}, \quad j = 1, \ldots, N.$$

From (5.23) and Lemma 5.5 we see that $z_j \in \Lambda$. Given $A \in P_0$, we have by Lemma 5.6 that the points $z_j$ with $j \in A$, belong to the same connected component of $\overline{\Lambda} \cup \text{supp} \mu_0$. Then there exists a connected set $\gamma_A \subset \overline{\Lambda} \cup \text{supp} \mu_0$ which is a finite union of bounded $C^1$ Jordan arcs satisfying

$$z_j \in \gamma_A, \quad \text{for every } j \in A.$$

For $L_{j,R}$ given by (5.21), define

$$\Gamma_A = \gamma_A \cup \bigcup_{j \in A} L_{j,R}.$$  (5.25)

Each $\Gamma_A$ is then a finite union of $C^1$ arcs. $\Gamma_A$ is connected and $\Gamma_A$ stretches out to infinity in each sector $S_j$ with $j \in A$. The union

$$\Gamma_0 = \bigcup_{A \in P_0} \Gamma_A$$

then satisfies all the requirements in Definition 2.2 and it follows that

$$\Gamma_0 \in \mathcal{T}(P).$$

Hence

$$I^\varphi(\Gamma_0) \leq \sup_{\Gamma \in \mathcal{T}(P)} I^\varphi(\Gamma).$$
Since $\Gamma_0 \subset \Gamma_0 \cup \text{supp } \mu_0$, we also have that
\begin{equation}
I^\varphi(\Gamma_0 \cup \text{supp } \mu_0) \leq I^\varphi(\Gamma_0).
\end{equation}

On the other hand, since $\Gamma_0 \subset \mathcal{X} \cup \text{supp } \mu_0$, the variational conditions (2.4), (2.5) are valid for $\Gamma_0 \cup \text{supp } \mu_0$, which means that $\mu_0$ is the equilibrium measure of $\Gamma_0 \cup \text{supp } \mu_0$ in the external field $\varphi$. Thus
\begin{equation}
I^\varphi(\Gamma_0 \cup \text{supp } \mu_0) = I^\varphi(\mu_0) = I^\varphi(F_0),
\end{equation}
since $\mu_0$ is the equilibrium measure in the external field of $F_0$ as well. Then we recall that $F_0$ is the maximizer of $I^\varphi$ over the class $\mathcal{F}_M$ which contains $\mathcal{T}_M$. Then by Proposition 5.1
\begin{equation}
I^\varphi(F_0) \geq \sup_{\Gamma \in \mathcal{T}_M} I^\varphi(\Gamma) = \sup_{\Gamma \in \mathcal{T}} I^\varphi(\Gamma)
\end{equation}
Combining the inequalities in (5.25)–(5.28), we finally get
\begin{equation}
I^\varphi(\Gamma_0) \leq \sup_{\Gamma \in \mathcal{T}} I^\varphi(\Gamma) \leq I^\varphi(F_0) = I^\varphi(\Gamma_0 \cup \text{supp } \mu_0) \leq I^\varphi(\Gamma_0)
\end{equation}
and so equality holds throughout in (5.29).

Let $\mu_1 = \mu(\varphi,F_0)$ be the equilibrium measure in the external field $\varphi$ of $\Gamma_0$. Then because of the equalities in (5.29), we have
\begin{equation}
I^\varphi(\mu_1) = I^\varphi(\Gamma_0) = I^\varphi(\Gamma_0 \cup \text{supp } \mu_0)
\end{equation}
Thus $\mu_1$ is a measure supported on $\Gamma_0 \subset \Gamma_0 \cup \text{supp } \mu_0$ whose energy in external field $\varphi$ is equal to the minimal energy for measures on $\Gamma_0 \cup \text{supp } \mu_0$. By uniqueness of equilibrium measure it follows that $\text{supp } \mu_0 \subset \Gamma_0$ and $\mu_0 = \mu_1$ is the equilibrium measure in the external field $\varphi$ of $\Gamma_0$.

Then $\Gamma_0 \in \mathcal{T}$ has the S-property in external field $\varphi$, see (5.18) in Proposition 5.4. Also the remaining statements of Theorem 2.3 follow from Proposition 5.4.

This completes the proof of Theorem 2.3, assuming the validity of Lemma 5.6.

5.4. Proof of Lemma 5.6. What remains is the proof of Lemma 5.6.

Proof of Lemma 5.6. We take $A \in \mathcal{P}_0$ and $j,k \in A$ with $j \neq k$ and our task is to prove that $\Lambda(j) = \Lambda(k)$.

In the proof we make use of the maximizer $F_0 \in \mathcal{F}_M$ and we know that $\mu_0$ is the equilibrium measure of $F_0$ in the external field $\varphi$. Thus by the variational conditions (2.4), (2.5),
\begin{equation}
F_0 \subset \{ z \in \mathbb{C} \mid U^{\mu_0}(z) + \frac{1}{2} \varphi(z) \geq I_0 \}.
\end{equation}

If $z \in F_0 \setminus \text{supp } \mu_0$ then $U^{\mu_0} + \frac{1}{2} \varphi$ is harmonic near $z$, and so cannot have a local maximum at $z$. It follows that
\begin{equation}
F_0 \subset \mathcal{X} \cup \text{supp } \mu_0.
\end{equation}
where we recall that $\Lambda$ is defined in (5.19).

Since $\mathcal{F}_M$ is the closure of $\mathcal{T}_M$, there exists a sequence $(\Gamma_n)_n$ in $\mathcal{T}_M$ such that
\begin{equation}
\Gamma_n \to F_0 \quad \text{in Hausdorff metric.}
\end{equation}
Let $n \in \mathbb{N}$. From item iii) of Definition 2.2 it follows that $\Gamma_n$ has a connected component that stretches out to infinity in the sectors $S_j$ and $S_k$. Then by dropping parts of that component that go to infinity in other sectors, and parts that make loops, we can find a subset $\tilde{\Gamma}_n \subset \Gamma_n$ that is a simple piecewise $C^1$ contour that
Figure 5. Pictorial example for the proof of Lemma 5.6 considering a polynomial $V$ of degree 5. The shaded region is $\Delta_M$. The black thick arcs belonging to a circle are the curves $\alpha_i$. The oriented contour is an example of $\tilde{\Gamma}_n$ for some $\Gamma_n$ in (5.31). The points $\xi_{i,n}, \eta_{i,n}$, $j = 1, \ldots, l$ are also represented in the figure, in this case with value $l = 3$.

goes from infinity in the sector $S_j$ to infinity in the sector $S_k$. Here simple means that it has no points of self-intersection. We consider $\tilde{\Gamma}_n$ as an oriented contour, and we talk about points on the contour that lie before or after other points. From the definition of $T_M$, see (5.4), we know that $\tilde{\Gamma}_n \cap \Delta_M = \emptyset$. (5.32)

The region $\Delta_M$ consists of $N$ connected components, all of them stretching out to infinity, as shown with the gray shaded regions in Figure 4 for the case $N = 5$.

Let $R$ be large enough such that the circle $|z| = R$ intersects each of the components of $\Delta_M$ along a circular arc. Then there are also $N$ circular subarcs of $|z| = R$ outside of $\Delta_M$, that we call $\alpha_1, \ldots, \alpha_N$ where $\alpha_i$ is in the direction of the sector $S_i$, for $i = 1, \ldots, N$. By taking $R$ large enough, we can also make sure that

$$\{z \in \alpha_i \mid U^{\mu_0}(z) + \frac{1}{2} \varphi(z) \geq l_0\} \subset \Lambda(i)$$

(5.33)

for each $i$, where $\Lambda(i)$ is as in (5.24).

We now follow the contour $\Gamma_n$ starting at infinity in the sector $S_j$. It will meet the arc $\alpha_j$. We let $\xi_{1,n}$ be the last point on $\tilde{\Gamma}_n$ that is on $\alpha_j$. Then $\tilde{\Gamma}_n$ enters into $D_R \setminus \Delta_M$ and never returns to $\alpha_j$. Since it ends up at infinity again, the contour has
to leave $D_R \setminus \Delta_M$ again, and it will do so along one of the arcs $\alpha_i$, with $i \neq j$. We let $\eta_{i,n}$ be the first point after $\xi_{1,n}$ which is on one of the $\alpha_i$ again, say $\eta_{1,n} \in \alpha_{j_1}$. If $j_1 = k$, then we stop. If $j_1 \neq k$, then we continue and we let $\xi_{2,n}$ be the last point on $\Gamma_n$ that is on $\alpha_{j_1}$. After $\xi_{2,n}$ the contour is going into $D_R \setminus \Delta_M$ again, since the contour has to end in the sector $S_k$ and $j_1 \neq k$. The contour will meet one of the $\alpha_i$ again, and we let $\eta_{2,n}$ be the first point after $\xi_{2,n}$ which is on one of the $\alpha_i$, say $\eta_{2,n} \in \alpha_{j_2}$. Then $j_2 \neq j$ and $j_2 \neq j_1$, since $\xi_{1,n}$ is the last point on $\alpha_j$ and $\xi_{2,n}$ is the last point on $\alpha_{j_1}$, and $\eta_{2,n}$ comes after these two points. If $j_2 = k$ then we stop, and otherwise we continue, see Figure 3 for an illustration.

Continuing this process, we find a sequence of distinct indices $j_0, j_1, \ldots, j_l$ for some $l \in \{1, \ldots, N-1\}$, where

$$j_0 = j, \quad j_l = k,$$

and points $\xi_{1,n}, \eta_{1,n}, \ldots, \xi_{l,n}, \eta_{l,n}$ on the contour $\Gamma_n$ with

$$\xi_{i,n} \in \alpha_{j_{i-1}}, \quad \eta_{i,n} \in \alpha_{j_i} \quad \text{for } i = 1, \ldots, l. \quad (5.34)$$

Furthermore, if we use $\Gamma_n(\xi, \eta)$ to denote the part of $\Gamma_n$ lying strictly between two points $\xi$ and $\eta$ of the contour, then

$$\Gamma_n(\xi_{i,n}, \eta_{i,n}) \subset D_R \setminus \Delta_M \quad \text{for } i = 1, \ldots, l. \quad (5.35)$$

The numbers $l, j_1, \ldots, j_{l-1}$ depend on the contour $\Gamma_n$ and so they depend on $n$. But these numbers are always positive integers smaller than or equal to $N$, so there are just finite possible choices for them. By passing to a subsequence of the sequence $(\Gamma_n)$, we may assume they are independent of $n$.

The set

$$\Gamma_n[\xi_{i,n}, \eta_{i,n}] = \Gamma_n(\xi_{i,n}, \eta_{i,n}) \cup \{\xi_{i,n}, \eta_{i,n}\}$$

is a connected closed subset of $\Gamma_n$ lying in the compact set $\overline{D_R \setminus \Delta_M}$. By compactness of the Hausdorff metric, there is a subsequence, say $N_1 \subset \mathbb{N}$, such that the limit

$$F_i = \lim_{n \to \infty, n \in N_1} \Gamma_n[\xi_{i,n}, \eta_{i,n}]. \quad (5.36)$$

eexists in the Hausdorff metric for every $i = 1, \ldots, l$.

The limit $F_i$ is connected, because this property is preserved by taking limits in the Hausdorff metric if all the sets involved are contained in a large, but fixed, compact set of $\mathbb{C}$, in this case $\overline{D_R}$, see (5.35). Also

$$F_i \subset F_0 \quad \text{for } i = 1, \ldots, l, \quad (5.37)$$

as follows from (5.31) and the fact $\Gamma_n[\xi_{i,n}, \eta_{i,n}] \subset \Gamma_n$.

By compactness of $\alpha_i$ we may also assume (by taking a further subsequence, if necessary) that the sequences $(\xi_{i,n})_n$ and $(\eta_{i,n})_n$ for $i = 1, \ldots, l$ converge along the same subsequence $N_1$, say

$$\xi_i = \lim_{n \to \infty, n \in N_1} \xi_{i,n}, \quad \eta_i = \lim_{n \to \infty, n \in N_1} \eta_{i,n}, \quad \text{for } i = 1, \ldots, l.$$

Then by (5.34) and (5.36)

$$\xi_i \in F_i \cap \alpha_{j_{i-1}}, \quad \eta_i \in F_i \cap \alpha_{j_i}, \quad (5.38)$$

for every $i = 1, \ldots, l$. From (5.38) and (5.33) it follows that

$$\xi_i \in \Lambda(j_{i-1}) \quad \eta_i \in \Lambda(j_i).$$
Each $F_i$ is a connected subset of $F_0$ and by (5.30) and (5.38) we see that $\xi_i$ and $\eta_i$ belong to the same connected component of $\Lambda \cap \text{supp} \mu_0$, which means that

$$\Lambda(j_{i-1}) = \Lambda(j_i) \quad \text{for } i = 1, \ldots, l.$$  

Thus $\Lambda(j_0) = \Lambda(j_l)$ which gives us $\Lambda(j) = \Lambda(k)$ as required. \hfill \Box

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