NATURE OF Hα SELECTED GALAXIES AT \( z > 2 \). I. MAIN-SEQUENCE AND DUSTY STAR-FORMING GALAXIES

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ABSTRACT

We present the results from our narrow-band imaging surveys of Hα emitters (HAEs) at \( z = 2.2 \) and \( z = 2.5 \) in the Subaru/XMM-Newton Deep survey Field with near-infrared camera MOIRCS on the Subaru Telescope. We have constructed a clean sample of 63 star-forming galaxies at \( z = 2.2 \) and 46 at \( z = 2.5 \). For 12 (or \( \sim 92\% \)) of 13 HAEs at \( z = 2.2 \), their Hα emission lines have been successfully detected by the spectroscopy. While about 42\% of the red, massive HAEs with \( M_* \geq 10^{10.8} M_\odot \) contain active galactic nuclei (AGNs), most of the blue, less massive ones are likely to be star-forming galaxies. This suggests that the AGN may play an important role in galaxy evolution at the late stage of truncation. For the HAEs excluding possible AGNs, we estimate the gas-phase metallicities on the basis of [N\text{II}]/Hα ratios, and find that the metallicities of the Hα selected galaxies at \( z = 2.2 \) are lower than those of local star-forming galaxies at fixed stellar mass, as shown by previous studies. Moreover, we present and discuss the so-called main sequence of star-forming galaxies at \( z > 2 \) based on our unique sample of HAEs. By correlating the level of dust extinction with the location on the main sequence, we find that there are two kinds/modes of dusty star-forming galaxies: starbursting galaxies and metal-rich normal star-forming galaxies.

Key words: galaxies: evolution – galaxies: high-redshift – galaxies: starburst

Online-only material: color figures

1. INTRODUCTION

The star formation rate (SFR) is one of the most fundamental parameters to characterize current activities of galaxies and its time evolution because new stars are continuously formed through cooling and contraction of cold gas in galaxies. The sketch of the cosmic evolution of star formation activities is known as the “Madau plot” (Madau et al. 1996). Hopkins & Beacom (2006) have compiled an extensive series of SFR measurements and illustrated the evolution of cosmic SFR density from \( z = 0 \) back to \( z \sim 6 \). The volume-averaged SFR (i.e., SFR density) increases by a factor of 10 from \( z = 0 \) to \( z \sim 1 \) and comes to its peak at \( z \sim 2 \), which indicates that a large fraction of stars in the present-day galaxies were formed beyond \( z \sim 1 \). The evolution of stellar mass density independently shows that about 50\% of stellar mass in the local galaxies are formed at \( z > 1 \) (e.g., Marchesini et al. 2009; Kajisawa et al. 2009). Therefore, star-forming galaxies at the peak epoch are the key population to understanding the formation and early evolution of galaxies.

While the SFR represents a differential rate of on-going star formation, the stellar mass (\( M_* \)) indicates the cumulative amount of past star formation activities. Therefore, these two are fundamental quantities to characterize star formation histories of galaxies. Daddi et al. (2007) find a tight correlation between stellar mass and SFR in galaxies at \( z \sim 2 \), called the “main sequence” of star-forming galaxies. A similar correlation is also found at \( z \sim 1 \) (Noeske et al. 2007) and in the local universe (Elbaz et al. 2007), while the cosmic SFR density declines significantly by more than an order of magnitude from \( z = 1 \) to 0.

In the SFR–\( M_* \) diagrams, some outliers have significantly higher specific SFRs (sSFRs) with respect to the main-sequence galaxies. Using the simplified model, Renzini (2009) has shown that a small difference in sSFR can be dramatically amplified in the subsequent evolution. While galaxies with high sSFRs undergo a rapid mass accretion, those with low sSFRs evolve with a moderate mass increase. In fact, recent observations of molecular gas in star-forming galaxies at \( z = 1–2 \) have revealed that they are very gas-rich systems, and suggest that two different star formation modes exist; a long-lasting mode for “normal” galaxies, and a short-lived intense mode for “starburst” galaxies (Daddi et al. 2010b).

Rodighiero et al. (2011) have demonstrated that the main-sequence galaxies represent 98\% of mass-selected star-forming galaxies and account for \( \sim 90\% \) of the cosmic SFR density at \( z \sim 2 \), with the sample detected by PACS 100 \( \mu \)m or 160 \( \mu \)m onboard Herschel as well as the optical/near-infrared (NIR) color-selected one. Some recent studies also suggest that the main sequence is independent of the environment in the local universe (Peng et al. 2010) and in the distant universe (Koyama et al. 2013b), supporting the fact that main-sequence galaxies represent a vast majority of star-forming galaxies at each redshift.

Despite its importance, the definition of the main sequence is largely dependent on a sample selection and the SFR indicator used. In this paper, we present the main sequence of star-forming galaxies at \( z > 2 \), which are selected from an Hα narrow-band (NB) imaging survey. The advantage of selecting galaxies with NB imaging is that we can construct a nearly SFR-limited, complete sample of star-forming galaxies which spans a broader range in stellar mass from \( M_* \sim 10^8 M_\odot \) to \( M_* > 10^11 M_\odot \). Also, the Hα line is one of the best SFR indicators that has many great advantages, such as being less affected by dust extinction, and having been well calibrated in the local universe (Hopkins et al. 2003). Recently, we have been conducting the “MAHALO–Subaru” project (MAApping HAlpha and Lines of Oxygen with Subaru; see an overview by Kodama
et al. 2013), which is a large, systematic survey of Hα emitters (HAEs) utilizing a unique set of NB filters. We target both biased high-density regions (Tanaka et al. 2011; Hayashi et al. 2012; Koyama et al. 2013a) and non-biased blank fields (Tadaki et al. 2011). In general fields, other large Hα surveys have been carried out at $z \sim 2$. Sobral et al. (2013) made a large sample of HAEs at $z = 2.23$ with a NB filter over $\sim 2$ deg$^2$ as a part of the High-$z$ Emission Line Survey (HiZELS) with Wide Field Camera (WFCAM) on UKIRT, and derived the Hα luminosity function and its evolution. Hayes et al. (2010) conducted the deepest survey of HAEs at this high redshift with HAWK-I on Very Large Telescope (VLT) over a 56 arcmin$^2$ area, and constrain the faint end slope of the luminosity function. Given such large or deep NB surveys providing statistical samples of star-forming galaxies at $z \sim 2$, our primary motivation is shifted from the stage of identifying high-$z$ star-forming galaxies to the next stage of understanding the physical processes in these galaxies that control their formation and early evolution.

This paper is structured as follows. In Section 2, we describe the existing data and our observations in Subaru/XMM-Newton Deep survey Field (SXDF; Furusawa et al. 2008). We show the target selections of HAEs in Section 3. In Section 4, their global properties such as SFRs, stellar masses, and metallicities are derived, and the mass–metallicity relation is presented. In Section 5, we discuss the relation between the main-sequence galaxies and dustiness of star formation activities and interpret the mass–metallicity relation observed at $z > 2$. We summarize our study in Section 6. Throughout this paper, we assume the cosmological parameters of $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_M = 0.3$, and $\Omega_{\Lambda} = 0.7$, and the Salpeter initial mass function (IMF) is adopted for the estimation of stellar masses and SFRs (Salpeter 1955).

2. DATA

We have conducted the HAE surveys in two redshift slices of $z = 2.2$ and $z = 2.5$ at a general field SXDF. Two NB filters, namely NB209 ($\lambda_c = 2.09$ $\mu$m, FWHM = 0.025 $\mu$m) and NB2315 ($\lambda_c = 2.315$ $\mu$m, FWHM = 0.027 $\mu$m), are used. Figure 1 shows the filter response functions of the two NB filters together with a broad-band (BB) filter $K$-band which measures the continuum level. The NB209 and NB2315 filters capture Hα emission lines from star-forming galaxies at $z = 2.191 \pm 0.020$ and $z = 2.525 \pm 0.021$, respectively. Note that $z \sim 2.5$ is the highest redshift where we can capture Hα lines with high sensitivity from ground-based telescopes.

2.1. Broad-band Data

In SXDF, there is an extensive data-set covering a wide wavelength from ultraviolet (UV) to mid-infrared. We use the publicly available multi-wavelength catalog (Galametz et al. 2013) from the Rainbow Database, which includes the $u$-band data from CFHT/Megacam (O. Almaini et al., in preparation), $B_-, V_-, R_-, i_-,$ and $z_-''$-band data from Subaru/Suprime-Cam (Furusawa et al. 2008), $Y_-$ and $K_s$-band data from VLT/HAWK-I (A. Fontana et al., in preparation), $J_-$, $H_-$, and $K$-bands data from UKIRT/WFCAM (UKIDSS Data Release 8; Lawrence et al. 2007), Spitzer/IRAC (SEDS; Ashby et al. 2013) and MIPS data (SpUDS; PI: J. Dunlop). Table 1 summarizes the existing BB data that are used in this paper, and their limiting magnitudes.

2.2. Narrow-band Imaging

The NB imaging observations were carried out with Multi-Object InfraRed Camera and Spectrograph (MORICS; Suzuki et al. 2008) on Subaru. MORICS is equipped with two Hawaii-2 detectors ($2048 \times 2048$), and provides a $7 \times 4'$ field of view per pointing. We made observations for seven nights in total, in 2010 October to November and 2011 September. The weather condition was fine during the observing runs and the seeing sizes were $0.4-0.7$ in FWHM. The point spread functions (PSFs) of all the images are eventually smoothed to $0.7$. Figure 2 shows the observed fields, and Table 2 lists the exposure times, seeing sizes (FWHM), field-of-views (FoVs), and the limiting magnitudes in the NB209 or NB2315 bands. The observed field is unique because the high-resolution optical/NIR images by ACS/WFC3 on the Hubble Space Telescope are both publicly available (Grogin et al. 2011). The field coverages are slightly different between the NB209 and NB2315 surveys. For NB2315, we spent four MORICS pointings which covers a contiguous field of 94 arcmin$^2$. For NB209, unfortunately, one of the two MORICS chips was unavailable at the time of observation due to a mechanical problem. Therefore, we spent eight MORICS pointings to cover a total area of 91 arcmin$^2$. These correspond to the co-moving volumes of $1.21 \times 10^3$ Mpc$^3$ and $1.26 \times 10^3$ Mpc$^3$ at $z = 2.2$ (NB209) and $z = 2.5$ (NB2315), respectively.

![Figure 1](https://example.com/figure1.png)

**Figure 1.** Transmission curves of NB209 and NB2315 filters on MOIRCS (solid line) and the WFCAM K-band filter (dashed line). The labels on the top axis indicate redshifts for HAEs.

**Table 1 Multi-wavelength Data in SXDF**

| Filter | Instrument | $m_{50,AB}$ | Reference |
|--------|------------|-------------|-----------|
| $u$    | CFHT/MegaCam | 27.68 | O. Almaini et al., in preparation |
| $B$    | Subaru/Suprime-Cam | 28.38 | Furusawa et al. (2008) |
| $V$    | Subaru/Suprime-Cam | 28.01 | Furusawa et al. (2008) |
| $R_c$  | Subaru/Suprime-Cam | 27.78 | Furusawa et al. (2008) |
| $i'$   | Subaru/Suprime-Cam | 27.69 | Furusawa et al. (2008) |
| $z'$   | Subaru/Suprime-Cam | 26.67 | Furusawa et al. (2008) |
| $Y$    | VLT/HAWK-I | 26.69 | A. Fontana et al., in preparation |
| $K_s$  | VLT/HAWK-I | 25.92 | A. Fontana et al., in preparation |
| $J$    | UKIRT/WFCAM | 25.63 | Lawrence et al. (2007) |
| $H$    | UKIRT/WFCAM | 24.76 | Lawrence et al. (2007) |
| $K$    | UKIRT/WFCAM | 25.39 | Lawrence et al. (2007) |
| 3.6 $\mu$m Spitzer/IRAC | 24.72 | Ashby et al. (2013) |
| 4.5 $\mu$m Spitzer/IRAC | 24.61 | Ashby et al. (2013) |
| 5.8 $\mu$m Spitzer/IRAC | 22.30 | SpUDS |
| 8.0 $\mu$m Spitzer/IRAC | 22.26 | SpUDS |
| 24 $\mu$m Spitzer/MIPS | 30–60 $\mu$Jy | SpUDS |
Figure 2. Field coverages of our NB surveys with MOIRCS are shown by green squares. Red squares and blue polygons indicate the areas covered by CANDELS WFC3 and ACS surveys, respectively.

(A color version of this figure is available in the online journal.)

Table 2
A Summary of our NB Imaging Observations with MOIRCS

| Field    | Exposure Time (min) | Seeing (") | FoV  |
|----------|---------------------|------------|------|
| NB209-F1 | 147                 | 0.5        | 4' x 3.5 | 23.5 |
| NB209-F1R| 140                 | 0.5        | 4' x 3.5 | 23.5 |
| NB209-F2 | 167                 | 0.6        | 4' x 3.5 | 23.6 |
| NB209-F2R| 167                 | 0.7        | 4' x 3.5 | 23.5 |
| NB209-F3 | 160                 | 0.7        | 4' x 3.5 | 23.6 |
| NB209-F3R| 127                 | 0.6        | 4' x 3.5 | 23.3 |
| NB209-F4 | 137                 | 0.6        | 4' x 3.5 | 23.6 |
| NB209-F4R| 177                 | 0.6        | 4' x 3.5 | 23.6 |
| NB2315-F1| 180                 | 0.6        | 4' x 7'  | 22.9 |
| NB2315-F2| 180                 | 0.6        | 4' x 7'  | 22.9 |
| NB2315-F3| 186                 | 0.6        | 4' x 7'  | 22.8 |
| NB2315-F4| 186                 | 0.6        | 4' x 7'  | 22.9 |

Data reduction is conducted by using the MOIRCS imaging pipeline software (MCSRED; Tanaka et al. 2011). Here, we briefly describe the procedures. First, a dome flat image is used to correct for any variabilities in sensitivity from pixel to pixel because a sky pattern remains in a self-sky flat image due to interferences of the OH sky emission lines in the case of NB imaging. The residual pattern, which is estimated by dividing a self-sky flat image by a dome flat image, is subtracted from each flat-fielded image. Then, a “median sky” image, which is made from some adjacent frames before and after the individual image to be processed, is also subtracted from the image. After each image is flat-fielded and sky-subtracted, a geometrical distortion is then corrected. Finally, all the processed images are combined with the inverse square of rms level of each frame as a weight. The photometric zero-point of our NB209/NB2315 images are determined using colors of unsaturated stars within the FoVs.

3. TARGET SELECTION

3.1. Narrow-band Emitters

We identify emission line galaxies on the basis of flux excesses in a NB as compared to a BB that covers the NB wavelength. First, the NB-detected catalogs are made from the images (H, K from WFCAM and NB from MOIRCS) using SExtractor (Bertin & Arnouts 1996). The pixel scales and PSF sizes of the BB images are matched to those of the NB images. The offset between these images is smaller than 0.05 arcsec. Source detections and photometries are carried out using the double image mode of the SExtractor. The extraction criterion is having at least nine pixels with fluxes above 1.5σ level in the NB images. The sources fainter than 5σ limiting magnitudes in the detection frame (NB209 or NB2315) are rejected. An aperture magnitude within a diameter of 1.6 is used to derive a color index. Note that these catalogs are used for the selection of NB emitters and the measurements of NB fluxes. For the following color selection (Figure 4) and the spectral energy distribution (SED) fitting (Section 4.1), the template-fitting photometries are used (see Galametz et al. 2013 for more details). Because the effective wavelengths of the NB209/NB2315 filters are slightly different from that of the K-band filter, we estimate the continuum level at the wavelengths of the NB209/NB2315 filters by taking a linear interpolation between H-band and...
Therefore, instead of using photometric redshifts, we define a technique is expected to be effective in discriminating HAEs at $z$ $\sim$ 2.2. As for the HAEs at $z$ $\sim$ 2.5, the $BiK$ and $iHK$ diagrams are used to reject $z$ $\sim$ 0.2 emitters and $z$ $\sim$ 3.6 emitters. The degrees of completeness and contamination in the color selections are estimated, based on the spectroscopically confirmed MODS sample. 98% (91%) of the MODS sample at $z$ $\sim$ 2.2 ($z$ $\sim$ 2.5) satisfy our color criteria and 5% (5%) of the MODS sample at other redshifts are included as contaminants. After these color selections, 63 HAEs at $z$ $\sim$ 2.2 and 46 at $z$ $\sim$ 2.5 are finally identified. While most of NB2315 emitters are HAEs at $z$ $\sim$ 2.5, the NB209 sample includes about 30 $[O\,\lambda\,\lambda\,\lambda]$ emitters at $z$ $\sim$ 3.2.

### 3.2. MOIRCS Spectroscopic Follow-up

In order to confirm that the HAEs, which will be selected in Section 3, are actually located at $z$ $\sim$ 2.2, and to measure their spectral properties such as the line ratio of $[N\,\lambda\,\lambda\,\lambda]/H\alpha$, we conducted a spectroscopic follow-up observation with MOIRCS on the multi-object spectroscopy (MOS) mode (Tokoku et al. 2006) in 2012 October to November. The HK500 low-resolution grizm, which covers the wavelength range of $\lambda$ $\sim$ 1.3–2.3 $\mu$m, was used. The slit width was 0.8 and its length was 10″–12″. The total on-source exposure time was 220 minutes.

Data reduction is conducted with the MOIRCS MOS Data Pipeline (MCSMDP; Yoshikawa et al. 2010) in the following way. First of all, bad pixels and cosmic-rays are removed. A sky frame is created in parallel to processing the object frames. The object frames and the sky frame are then divided by the dome flat for flat-fielding. Distortion is corrected by using the database for the imaging mode of MCSRED. The wavelength calibration is performed by using the OH lines in the sky frame. After a sky pattern is roughly removed by subtracting the adjacent dithered image (A–B subtraction) for each object frame, the residual background emission is subtracted by using the task “background” of IRAF. Finally, all frames are combined with weights in exposure times. To calibrate the telluric absorption and the difference in sensitivity at each wavelength, the object frames are divided by the spectrum of a spectroscopic standard.
star (A1-type) and multiplied by the model spectrum (Castelli & Kurucz 2004). For a slit spectroscopy, a slit loss must be properly corrected for in order to estimate a total flux. However, the slit loss strongly depends on the size of a galaxy. Therefore, in this work, we do not use the absolute flux derived from the spectroscopy and discuss only the redshifts and the line ratios of the HAEs.

For the target selection of spectroscopy, we gave a high priority on objects with a bright Hα emission measured from the NB data. With the MOIRCS, the NIR spectra of 13 HAEs at $z = 2.2$ have been obtained. We successfully detected Hα and [N II] emission lines above 3σ from 12 and 8 objects, respectively. Table 3 lists their spectroscopic redshifts, and Figure 5 shows their spectra. For only one object, no line is seen in its spectrum though the expected Hα flux density from our NB data is large enough to be detected. The non detection might be due to a large flux loss by the slit because it is a diffuse source. In order to measure the spectroscopic redshifts and line ratios, the reduced spectra are fit by a three-component (Hα and [N II]λ6548, λ6583) Gaussian model, with the free parameters of redshift, line width, and flux density of each line. The continuum is first subtracted off by fitting the spectra with a linear function over the wavelength range of 2.08–2.1 μm. The line ratio of [N II]λ6583/[N II]λ6548 is fixed to 3.0 (Storey & Zeippen 2000). The sky noise is estimated from a Poisson error of the sky frame without a background subtraction.

For SXDF-NB209-4, 34, 43, and 60, [N II] lines are not detected. SXDF-NB209-34 has been confirmed to be located at $z = 2.2$ as its [O III] and Hβ emission lines are detected. Although we do not yet have conclusive evidence at this stage that the detected lines are Hα for SXDF-NB209-4, 43, and 60, our spectroscopic follow-up observations do support that our criteria of identifying HAEs based on NB and BB data should work very well. We hereafter regard all of the 12 sources with line detections as HAEs.

4. GALAXY PROPERTIES
4.1. SED Fitting

Multi BB photometries provide SEDs of individual galaxies, which give us valuable information of physical properties such as stellar mass, age, and amount of dust extinction. We fit the photometric SEDs traced by 12 bands ($u, B, V, R, i', z', Y, J, H, K_s, 3.6 \mu m, 4.5 \mu m$) with the stellar population synthesis model of Bruzual & Charlot (2003). We adopt the solar metallicity. The SED fitting with a standard χ² minimization procedure is applied by using hyperz (Bolzonella et al. 2000). The redshift is fixed to $z = 2.19$ and $z = 2.53$ for the NB209 and NB2315 HAEs, respectively. We use the libraries with exponentially declining star formation histories: SFR∝$e^{−t/\tau}$ with $\tau = 0.01, 0.02, 0.05, 0.10, 0.20, 0.50, 1.0, 2.0, 5.0$ Gyr. Salpeter IMF is adopted for the estimation of stellar masses and SFRs. The dust attenuation law for starburst galaxies (Calzetti et al. 2000) is applied to the model spectra. The best fitting template gives stellar mass, age, $A_V$, and $\tau$. Although the estimated age and the amount of dust extinction by SED...
fitting are not very reliable due to a degeneracy, the stellar mass is a rather robust quantity if IMF is given. Ideally, we should accurately measure the amount of dust extinction by the Balmer decrement method. However, Hβ lines are generally too weak to be detected for high-z galaxies, and we ought to rely on the values of dust extinction derived from the SED fitting.

4.2. AGN Contribution

Ideally, we can distinguish between star-forming galaxies and active galactic nuclei (AGNs) by measuring emission line ratios of [O iii]/Hβ and [N ii]/Hα and plot them on the so-called BPT diagram (Baldwin et al. 1981). We detected both Hβ and [O iii] emission lines for only one target (SXDF-NB209-34). For this object, the line ratios of log ([O iii]/Hβ) = 0.73 and log ([N ii]/Hα) = -0.96 indicate that this is a star-forming galaxy. Unfortunately, neither Hβ nor [O iii] lines are detected for the other objects, and so we are obliged to make a diagnosis based solely on the line ratio of [N ii]/Hα. [N ii]/Hα is particularly sensitive to shock excitation or the presence of a hard ionizing radiation field around an AGN. Kewley & Dopita (2002) put a theoretical upper limit on [N ii]/Hα ratio based on photo ionization and stellar population synthesis models. For star-forming galaxies, the ratio of [N ii]/Hα should always be less than 0.6. On the BPT diagram, Seyfert galaxies are often defined to have [O iii]/Hβ > 3 and [N ii]/Hα > 0.6, and low-ionization nuclear emission-line regions to have [O iii]/Hβ < 3 and [N ii]/Hα > 0.6 (Kauffmann et al. 2003a). Therefore, we regard the three HAEs (SXDF-NB209-2, 9, and 16) with [N ii]/Hα > 0.6 as AGNs.

One HAE (SXDF-NB209-1) has been diagnosed as a narrow-line AGN on the basis of high-ionization ultraviolet emission lines such as C iv and [Ne v] by the previous study (Simpson et al. 2012). The diagnostic with IRAC four-band photometries is also useful for identifying AGNs. The mid-infrared SEDs of AGNs are well fitted by a power law because the thermal emission by hot dust dominates the SEDs of host galaxies (Donley et al. 2007). Donley et al. (2008) find that a power-law selection of α < −0.5 recovers the majority of high-quality AGN candidates. While the IRAC SEDs of most our HAEs are

![Figure 5. Hα line spectra of our 12 NB209 HAEs. Red lines show the observed spectra. Vertical dotted lines indicate the locations of Hα and [N II]λ6548, λ6583 lines. Gray zones show the sky noise levels. Flux is calibrated by using a star in the 2MASS catalog (Skrutskie et al. 2006), although we do not rely on the flux calibration in this paper.](image)
not fitted by a power law, one HAE (SXDF-NB209-42) satisfied the criterion of $\alpha < -0.5$, suggesting that it hosts an AGN. In total, we have identified five AGN candidates. In SXDF, deep X-ray data by XMM-Newton are available (Ueda et al. 2008), which enable us to identify AGNs. However, none of our five AGN candidates are detected in X-ray, suggesting that they are obscured AGN by dust and gas.

Figure 6 shows the color–magnitude diagram for HAEs at $z = 2.2$. We also plot seven HAEs at $z = 2.2$ in GOODS-N field, which are identified by our previous NB209 survey on MOIRCS and confirmed by the follow-up spectroscopy (Tadaki et al. 2011). One of them is an X-ray-detected AGN. Galaxies are known to show a color bimodality; blue cloud galaxies and red sequence galaxies (Kauffmann et al. 2003b). The HAEs are naturally supposed to be blue star-forming galaxies. However, some of these red candidates are detected in X-ray, suggesting that they are obscured AGN by dust and gas.

4.3. Star Formation Rates

It is standard for us to estimate SFRs of our HAEs with the following standard calibrations of Kennicutt (1998);

$\text{SFR}_{\text{H}}[M_\odot \text{yr}^{-1}] = 7.9 \times 10^{-42}L_{\text{H}}[\text{erg s}^{-1}], \tag{4}$

$\text{SFR}_{\text{UV}}[M_\odot \text{yr}^{-1}] = 1.4 \times 10^{-28}L_{\text{UV}}[\text{erg s}^{-1} \text{Hz}^{-1}], \tag{5}$

where SFR$_{\text{H}}$ and SFR$_{\text{UV}}$ indicate dust-uncorrected SFRs based on $\text{H}_{\alpha}$ and rest-frame UV luminosities, respectively. Since the composite UV spectrum is nearly flat over the wavelength range of 1500–2800 Å, the UV luminosity is calculated by linearly interpolating the $B$-, $V$-, and $R_c$-band luminosities as

$\text{m}_{1500} = 0.7B + 0.3V$ for HAEs at $z = 2.2$, \tag{6}

$\text{m}_{1500} = 1.2V - 0.2R_c$ for HAEs at $z = 2.5$, \tag{7}

where $m_{1500}$ is the interpolated magnitude at the rest-frame wavelength of 1500 Å.

The HAE surveys with NB filters allow us to measure the $\text{H}_\alpha$ fluxes from the flux densities in NB and BB, hence SFRs for all galaxies (Koyama et al. 2010). The NB flux density can be defined as $f_{\text{SB}} = f_c + F_{\text{line}}/\Delta_{\text{SB}}$, where $f_c$ is the continuum flux density, $F_{\text{line}}$ is the emission-line flux, and $\Delta$ denotes FWHMs of the filters. The BB flux density is also defined as...
forming galaxies, is conveniently parameterized as

\[ SFR = \alpha \times M_\star^\beta. \]  

Figure 8. Left: Hα-based SFRs (SFRH\text{\textsc{\textsc{\textalpha}}}) vs. UV-based SFRs (SFRUV) without dust-correction for the combined sample of HAEs at \( z = 2.2 \) and \( z = 2.5 \). Right: same as the left panel, but with the dust extinction correction based on the SED fitting.

(A color version of this figure is available in the online journal.)

\[ f_{BB} = f_c + F_{line}/\Delta BB. \] Therefore, the line flux, continuum flux density, and equivalent width (EW) are calculated as

\[ F_{line} = \Delta_{NB} f_{NB} f_{BB}/(1 - \Delta_{NB}/\Delta_{BB}), \]  

\[ f_c = f_{NB} f_{BB}/(1 - \Delta_{NB}/\Delta_{BB}), \]  

\[ EW_{\text{rest}} = F_{line}/f_c(1 + z)^{-1}. \]  

Note that a [N\text{\textsc{\textsc{\textalpha}}}] line also contributes to a NB209/NB2315 flux. Sobral et al. (2012) estimated the [N\text{\textsc{\textsc{\textalpha}}}] contamination for the full Sloan Digital Sky Survey (SDSS) sample as a function of total rest-frame EW and compute the polynomial approximation of the relationship between log([N\text{\textsc{\textsc{\textalpha}}}]/Hα) (denoted as \( f \)) and log(\( EW_{\text{rest}}([N\text{\textsc{\textsc{\textalpha}}}]+H\alpha] \)) (denoted as \( E \)), which is presented as

\[ f = -0.924 + 4.802 E - 8.892 E^2 + 6.701 E^3 - 2.27E^4 + 0.279E^5. \]  

We use this polynomial function to correct for a [N\text{\textsc{\textsc{\textalpha}}}] contamination to derive a pure Hα flux. The average correction is [N\text{\textsc{\textsc{\textalpha}}}]/Hα = 0.16.

Figure 8 shows the SFRH\text{\textsc{\textalpha}} versus SFRUV diagram for the HAEs at \( z = 2.2 \) and 2.5 with and without dust extinction correction (right and left panels, respectively). The dust uncorrected SFRUV is not correlated to SFRH\text{\textsc{\textalpha}} (left panel). We use the amount of dust extinction derived from the SED fitting in Section 4.1 and assume equal extinction between UV (stellar) and Hα (nebular) rather than \( E(B-V)_{\text{stellar}} = 0.44 E(B-V)_{\text{nebular}} \) (Calzetti et al. 2000), because the SFRH\text{\textsc{\textalpha}} is significantly overestimated with respect to the SFRUV in the latter situation (Erb et al. 2006b). Once we apply such corrections, the right panel of Figure 8 shows a good correlation between the two measurements of SFRs from two independent indicators. Hereafter, we use dust-corrected SFRs based on Hα luminosities, in order to define the main sequence of star-forming galaxies at \( z > 2 \) presented in Section 4.4.

4.4. Main Sequence of Star-forming Galaxies

The \( M_\star - SFR \) relation, named as the main sequence of star-forming galaxies, is conveniently parameterized as

\[ SFR = \alpha \times M_\star^\beta. \]  

Based on the fact that the majority of star-forming galaxies are located on this sequence, it must have essential information on how star formation in galaxies are generally regulated, and it must be playing a vital role in galaxy evolution. In spite of its vital importance, the values of \( \alpha \) and \( \beta \) are still quite uncertain especially at high redshifts. This is because the results are highly dependent both on the sample selection and on the SFR indicator that is used.

Due to our relatively unbiased sample of star-forming galaxies and the robust measure of SFRs based on Hα, we are able to investigate such a relationship for star-forming galaxies at \( z = 2.2 \) and 2.5. Figure 9 shows the SFR–M\text{\textsc{\textalpha}} diagram for the HAEs. The star-forming activities of HAEs at \( z = 2.5 \) seem to be systematically higher than those at \( z = 2.2 \) at the fixed stellar mass. It is not clear at this moment whether such difference is due to an intrinsic evolution of galaxies with time, or an environmental dependence (cosmic variance), or simply due to small number statistics. We will investigate this potentially interesting result in greater detail in a future paper. In this paper, we do not consider this possible difference further, and we fit a relation for the combined sample of \( z = 2.2 \) and 2.5 with a stellar mass cut of \( M_\star > 10^{10.5}M_\odot \) since our HAE sample becomes significantly incomplete below it. Also, we do not use massive galaxies with \( M_\star > 10^{11.5}M_\odot \), in which the AGN fraction could increase (Section 4.2), for the fit. The best fit line is SFR = 238 M\text{\textsc{\textalpha}}^{0.94}, where M11 shows the stellar mass in the unit of \( 10^{11}M_\odot \). The best fit line is SFR = 200 M\text{\textalpha}^{0.53} as the main sequence of color-selected star-forming galaxies at \( z = 2 \). The best fitted main sequence is well suited to that defined by Daddi et al. (2007) although they estimated SFRs by using UV and IR luminosities rather than Hα emission lines, and the sample...
the mass–metallicity relation with the stacked spectra of 87 rest-
Hayashi et al. (2009) and Onodera et al. (2010) report that
that their metallicities are systematically lower than those of
objects with a large line flux. For the objects whose \([\text{N} \text{II}]\)
Note that our targets of spectroscopic observations are biased to
investigate the mass–metallicity relation further with a larger,
to the difference in the sample selection, it is important to
much higher metallicities which are almost comparable to those
optical–NIR color-selected star-forming galaxies at
formation histories. If an \(H\alpha\) emission line is free from AGN
emission line is much less attenuated by\) dust extinction than the UV light. Most of the MIPS-detected
HAEs actually show high values of dustiness (>5). We could
in principle estimate the amount of dust extinction \((A_V)\) directly
from the dustiness index by assuming the extinction curve such as
Calzetti et al. (2000): \(A_V = [2.5 \log(SFR_{H\alpha}/SFR_{UV}) + 0.4]/1.07\). However, we choose to use the dustiness index
because of the uncertainties in the extinction law.
In Figure 11, we find that the HAEs at the massive end of the main sequence tend to be detected at MIPS 24 \(\mu m\), indicating that they are dusty star-forming galaxies. However, for the galaxies below the detection limit of the MIPS data, we cannot evaluate how dusty they are from the MIPS data alone. To investigate a possible difference in dusty nature between the normal mode of star formation on the main sequence and the burst mode of star formation above the main sequence, we define the “dustiness” index with the following equation,

\[
\text{dustiness} = \frac{SFR_{H\alpha\, \text{dust - uncorrected}}}{SFR_{UV\, \text{dust - uncorrected}}}.
\]

For dusty star-forming galaxies, the dustiness index should be high because the \(H\alpha\) emission line is much less attenuated by dust extinction than the UV light. Most of the MIPS-detected HAEs actually show high values of dustiness (>5). We could in principle estimate the amount of dust extinction \((A_V)\) directly from the dustiness index by assuming the extinction curve such as Calzetti et al. (2000): \(A_V = [2.5 \log(SFR_{H\alpha}/SFR_{UV}) + 0.4]/1.07\). However, we choose to use the dustiness index because of the uncertainties in the extinction law.

Figure 11 plots the sSFRs as a function of stellar mass for our HAEs. It seems that the dustiness depends on two
parameters: offset from the main sequence in SFR, \(\triangle\log(\text{SFR}) = \log(\text{SFR}) - \log(\text{SFR}_{\text{MS}})\), and stellar mass. The HAEs with enhanced SFRs relative to the main-sequence galaxies tend to have dusty star formation. On the other hand, the dustiness is also clearly high at the massive end despite declined star formation activities (low sSFRs). These two interesting trends have been pointed out also by Wuyts et al. (2011), who find that \(\text{SFR}_{\text{R}}/\text{SFR}_{\text{UV}}\) is correlated with both star formation activities and stellar mass. We investigate the dependence of dustiness on two parameters separately by fixing one parameter and allowing the other parameter to change, and vice versa (Figure 12). For galaxies with \(M_\star < 10^{10.5}\, M_\odot\), the dustiness is closely correlated with \(\Delta \log(\text{SFR})\). This result may give us a hint on why the star formation activity is boosted in the galaxies with large \(\Delta \log(\text{SFR})\). In the case of clump migration (Ceverino et al. 2010; Inoue & Saitoh 2012) and gas-rich mergers, star-forming regions are likely to be concentrated toward galaxy centers due to an inflow of gas that has lost its angular momentum. Such compact dusty star-forming regions would lead to a large dust attenuation. We may therefore call such origin of dusty galaxies “starburst mode.”
Table 4

| ID     | R.A. (J2000) | Decl. (J2000) | z_Hα | M_∗ \(\times10^{10}M_\odot\) | [N II]/Hα | 12+log[O/H] |
|--------|--------------|--------------|------|-------------------------------|------------|-------------|
| MODS-487 | 12 35 58.32 | 62 11 55.7 | 2.176 | 1.7 | <0.33 | <8.63 |
| MODS-7651 | 12 37 01.97 | 62 15 50.0 | 2.188 | 0.9 | <0.07 | <8.24 |
| MODS-7773 | 12 36 56.74 | 62 16 43.7 | 2.190 | 5.9 | <0.26 | <8.57 |
| MODS-7889 | 12 36 52.90 | 62 17 26.5 | 2.187 | 0.5 | <0.09 | <8.30 |
| MODS-9441 | 12 37 18.22 | 62 16 53.4 | 2.196 | 1.3 | <0.07 | <8.24 |
| MODS-10522 | 12 37 25.54 | 62 19 09.8 | 2.192 | 0.7 | <0.13 | <8.39 |

Figure 11. Specific star formation rates vs. stellar masses for the HAEs. Color scale indicates the dustiness \((=SFR_{H\alpha}/SFR_{UV})\). The red line shows the main sequence of HAEs at \(z>2\). Magenta squares indicate the MIPS 24 \(\mu\)m detected sources.

(A color version of this figure is available in the online journal.)

On the other hand, the dustiness seems to be constantly high regardless of the level of star formation activities for the massive galaxies with \(M_\ast > 10^{10.5}M_\odot\). This suggests that stellar mass also contributes to the dust obscuration. In fact, at fixed \(\Delta\log(\text{SFR})\), the dustiness seems to increase along the main sequence from low mass to high mass ends. Such dependence indicates an additional process to increase dust extinction as stellar mass is increased. Star-forming galaxies continue to grow in stellar mass and move up along the main sequence as more and more new stars are formed. During the growth, chemical evolution must progress in the galaxies and become progressively more metal rich. Since the dust production is strongly correlated with metallicity, galaxies would become dustier as a result. Therefore, \(H\alpha\) emission lines and UV photons radiated from the ionized gas and massive stars, are strongly attenuated by dust, and they would appear as massive, dusty galaxies. In fact, in the local universe, it is found that the amount of dust extinction roughly correlates with stellar mass (Garn & Best 2010; Gilbank et al. 2010; Sobral et al. 2012). We may call such origin of dusty galaxies “metal-enrichment mode.”

In summary, we propose that the two modes/origins of dusty star-forming galaxies at \(z \sim 2\), are the starburst mode and the metal-enrichment mode.

5.2. Origin of the Metallicity Variation at Fixed Stellar Mass

We find that the metallicities of HAEs at \(z > 2\) are significantly lower compared to local galaxies at the fixed stellar mass (Section 4.5). It is not clear, however, why galaxies with the same stellar mass show such variations in metallicity despite the fact that they have produced the same amount of stars in the past. Recently, a fundamental relationship among three parameters, namely, mass, SFR, and metallicity, was presented, which can account for the source of scatter and the cosmic evolution of the mass–metallicity relation (e.g., Lara-López et al. 2010; Niino 2012). Mannucci et al. (2010) find that the metallicity is also correlated with SFR at fixed mass, and the scatter is significantly reduced by introducing SFR as a second parameter. Similarly, Yabe et al. (2012) demonstrate that galaxies with
large SFR have lower metallicity, with 71 star-forming galaxies at $z \sim 1.4$.

In order to investigate the effect of remaining molecular gas in the mass–metallicity relation, we present the mass–metallicity relation for the SDSS sample at $0.025 < z < 0.05$ whose CO emission lines are detected in Figure 13. Here we use the spectroscopic SDSS data (DR7; Abazajian et al. 2009) and the catalog of CO Legacy Data base for the GASS (GALEX Arecibo Sloan Digital Sky Survey) survey (COLD GASS; Saintonge et al. 2011). Molecular gas masses are derived from CO luminosities, with the conversion factor of $4.35 M_\odot / (K \text{ km s}^{-1} \text{pc}^{-2})$ for normal star-forming galaxies and $1.0 M_\odot / (K \text{ km s}^{-1} \text{pc}^{-2})$ in systems with high luminosities ($L_{\text{IR}} > 10^{11} L_\odot$) and warm dust temperatures ($S_{850 \mu \text{m}} / S_{1000 \mu \text{m}} > 0.5$) (Saintonge et al. 2012). The metallicities of galaxies at a given stellar mass show some scatter, but there is a clear trend that the galaxies with a larger gas mass fraction, $M_{\text{gas}} / (M_\ast + M_{\text{gas}})$, tend to have lower metallicities, indicating that the gaseous metallicity is well correlated to gas mass and its fraction at fixed stellar mass in the local universe. Bothwell et al. (2013) also find a dependence of metallicity on H$\alpha$ gas mass, supporting our result. Galaxies with a large gas reservoir have a potential to evolve into more massive and metal-rich galaxies by subsequent star formation, which would then be eventually located on the more massive side of the local mass–metallicity relation. While gas fractions of local star-forming galaxies are distributed in a relatively narrow range (5%–15%), high-$z$ galaxies are more gas rich (Tacconi et al. 2010; Daddi et al. 2010a). Therefore, the dependence of metallicity on the gas mass fraction would become more significant and visible. In order to reveal the origin of the metallicity variation, we need to measure the amount of remaining cold gas and its fraction for a statistical sample of star-forming galaxies at $z > 2$ with, e.g., ALMA. We can then confirm the hypothesis that the metallicity variation at a given stellar mass is due to the scatter in the gas mass fraction.

6. SUMMARY

We have conducted NB imaging and spectroscopic surveys of HAEs at $z = 2.2$ and $z = 2.5$ in SXDF, using MOIRCS on Subaru Telescope. Our survey has identified 109 HAEs at $z > 2$ in total over a total of 180 arcmin$^2$ area. We have confirmed probable H$\alpha$ emission lines for 12 out of 13 targets by the spectroscopic follow-up observations. Therefore, our technique of searching for HAEs based on the excess fluxes in NB and multi-color selection, is proven to be a robust and efficient method. Based on this unique, clean sample of star-forming galaxies at $z \sim 2$, we have investigated their global properties in this paper.

1. The diagnostics based on the [N II]/H$\alpha$ line ratio, high-ionization ultraviolet emission lines, or X-ray detections, shows that about 42% of the red, massive HAEs with $M_\ast > 10^{10.5} M_\odot$ contain AGNs, implying that the AGN feedback may be contributing to quenching star formation in these massive systems. To further investigate the effects of AGNs on the evolution of star-forming galaxies, we need to make systematic, extensive spectroscopic observations over a large sample that spans a wide range in stellar mass.

2. The HAEs at $z = 2.2$ and $z = 2.5$ exhibit a well-defined main sequence represented by $\text{SFR} = 238 M_\odot$ (Saintonge et al. 2011). We find that the dustiness index ($\text{SFR}_\text{IR}/\text{SFR}_\text{UV}$) of star-forming galaxies is dependent on two parameters: offset from the main sequence, $\Delta \log(\text{SFR})$, and the stellar mass. Galaxies with high SFRs with respect to the main sequence tend to have high dustiness. On the other hand, massive star-forming galaxies also tend to be dusty, probably because they are more metal rich and contain larger amounts of dust. Although it is widely recognized that dusty galaxies such as SMGs are merger-driven starburst galaxies (Engel et al. 2010), our result suggests that some dusty HAEs are not necessarily indicative of such populations and they can be more like metal-rich normal star-forming galaxies.

3. The metallicities of HAEs are roughly consistent with the typical values of the UV-selected galaxies at $z \sim 2.2$ (Erb et al. 2006a). Massive HAEs at $z > 2$ have already evolved into a metal-rich system to the same level as the local star-forming galaxies, while the metallicities of less massive galaxies are systematically lower, and their dispersion is large. The SDSS sample at $z \sim 0.1$ with CO detections suggests that the scatter in the mass–metallicity relation originated from the different amounts of molecular gas fraction. Gas-rich galaxies tend to have lower metallicities at fixed stellar mass compared to gas poor ones. Therefore, at $z > 2$ when the gas mass fraction of galaxies is larger (~40%) and the scatter would also be large, the metallicities at a given stellar mass would naturally show a large scatter as observed. The measurements of cold gas with ALMA and JVLA would reveal the dependence of metallicities on gas mass fraction at $z > 2$ directly for the first time, and this may lead us to establish a new fundamental metallicity relation among three parameters, namely, stellar mass, metallicity, and gas mass fraction.

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ERRATUM: “NATURE OF Hα SELECTED GALAXIES AT z > 2. I. MAIN-SEQUENCE AND DUSTY STAR-FORMING GALAXIES” (2013, ApJ, 778, 114)

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An error has been detected in the calculation of the line flux for Hα emitters (HAEs) at z = 2.5. This leads to the overestimation of ~50% in the line flux, and the number of the final sample at z = 2.5 decreases from 46 to 44. Figure 9 shows the modified diagram of the star formation rate (SFR) versus the stellar mass. We refit a main sequence to the HAEs at z = 2.2 and 2.5 and derived the best-fit line of SFR = 132 $M_\odot$ yr$^{-1}$. This revision, however, has little impact on the conclusions of the paper.

We also found a typo in Equation (9), and it is corrected as follows:

$$f_c = \frac{f_{BB} - f_{NB}(\Delta_{NB}/\Delta_{BB})}{1 - \Delta_{NB}/\Delta_{BB}}.$$  (9)

![Graph showing modified star formation rates of HAEs at z = 2.2 and z = 2.5](image)

**Figure 9.** Star formation rates of the combined sample of HAEs at z = 2.2 (green) and z = 2.5 (blue) plotted against stellar masses. We use the Hα-based SFRs with the dust extinction correction. The errors in the SFRs are estimated from the photometric errors ($\sigma = 1$). The red solid line indicates the best-fit line to our data points (SFR = 132 $M_\odot$ yr$^{-1}$). The dashed line presents the main sequence of star-forming galaxies at z ~ 2 defined by Daddi et al. (2007). Magenta squares indicate the MIPS 24 μm detected sources.

(A color version of this figure is available in the online journal.)

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