Long-path averaged mixing ratios of O$_3$ and NO$_2$ in the free troposphere from mountain MAX-DOAS

L. Gomez$^{1,2}$, M. Navarro-Comas$^1$, O. Puentedura$^1$, Y. Gonzalez$^2$, E. Cuevas$^3$, and M. Gil-Ojeda$^1$

$^1$Instituto Nacional de Técnica Aeroespacial (INTA), Área de Investigación e Instrumentación Atmosférica, Ctra Ajalvir km4, 28850, Torrejón de Ardoz, Madrid, Spain
$^2$Groupe de Spectrométrie Moléculaire et Atmosphérique, URM CNRS 7331, UFR Sciences Exactes et Naturelles, Moulin de la Housse, BP 1039, 51687 Reims CEDEX 2, France
$^3$Centro de Investigación Atmosférica de Izaña (Agencia Estatal de Meteorología –AEMET), Santa Cruz de Tenerife, Spain

Correspondence to: L. Gomez (lauragm.madrid@gmail.com)

Received: 7 May 2013 – Published in Atmos. Meas. Tech. Discuss.: 5 September 2013
Revised: 28 August 2014 – Accepted: 29 August 2014 – Published: 7 October 2014

Abstract. A new approximation is proposed to estimate O$_3$ and NO$_2$ mixing ratios in the northern subtropical free troposphere (FT). The proposed method uses O$_4$ slant column densities (SCDs) at horizontal and near-zenith geometries to estimate a station-level differential path. The modified geometrical approach (MGA) is a simple method that takes advantage of a very long horizontal path to retrieve mixing ratios in the range of a few pptv. The methodology is presented, and the possible limitations are discussed. Multi-axis differential optical absorption spectroscopy (MAX-DOAS) high-mountain measurements recorded at the Izaña observatory (28°18’ N, 16°29’ W) are used in this study. The results show that under low aerosol loading, O$_3$ and NO$_2$ mixing ratios can be retrieved even at very low concentrations. The obtained mixing ratios are compared with those provided by in situ instrumentation at the observatory. The MGA reproduces the O$_3$ mixing ratio measured by the in situ instrumentation with a difference of 28%. The different air masses scanned by each instrument are identified as a cause of the discrepancy between the O$_3$ observed by MAX-DOAS and the in situ measurements. The NO$_2$ is in the range of 20–40 ppt, which is below the detection limit of the in situ instrumentation, but it is in agreement with measurements from previous studies for similar conditions.

1 Introduction

The knowledge of reactive trace gas distributions in populated areas around the world has greatly improved recently due to the extensive deployment of measurement networks on local to regional scales. Field instrumentation generally provides information on the surface concentrations, rather than the vertical distributions, of the measured species. As a consequence, knowledge of free troposphere (FT) tracer distributions is lacking. Recently, remote sensing instrumentation based on spectroscopic analysis Multi-axis differential optical absorption spectroscopy (MAX-DOAS) has demonstrated success in obtaining profiles in polluted areas and has been used extensively for pollution studies (i.e. Baidar et al., 2013; Hendrick et al., 2014) and satellite validation purposes (Peters et al., 2012; Ma et al., 2013). The use of MAX-DOAS has been essentially restricted to the boundary layer (BL) where most of the pollutants are located. Under stable conditions, a BL thermal inversion limits the upward ventilation of pollutants and other surface-produced tracers; as a result, the FT concentrations of gases such as NO$_2$ are much lower (possibly by orders of magnitude). Occasionally, due to favourable weather conditions or to mechanical forcing on given orographies, polluted air masses may be injected into the FT (Thakur et al., 1999).

The knowledge of tracers in the FT is relevant to determine atmospheric reference background conditions for pollution studies (Engardt et al., 2009; Beelen et al., 2009) and to study chemical processes occurring during long-range
transport (Thakur et al., 1999). Pollutants can be transported over long distances through the FT where fast winds exist and the lifetime of minor species is longer than that in the BL (Stohl and Trickl, 1999; Stohl et al., 2003; Liang et al., 2004). Unfortunately, little is known of the amounts of BL active species in the FT that are anthropogenically produced (Wenig et al., 2003); the downward exchange between the FT and the lower troposphere is also unclear (Zyryanov et al., 2012).

Routine measurements of background levels of chemically active trace gases in the FT are very scarce. FT measurements have been recorded during limited time periods on airborne platforms with high-quality instrumentation (Robinson et al., 2005; Mao et al., 2006; Martin et al., 2006; Oltmans et al., 1996; Thompson et al., 2011; Baidar et al., 2013), but the high cost of the flights prevents their use for long-term studies. Continuous monitoring programmes have been limited to the use of in situ instruments located in the few existing mountain observatories around the world. However, the measurements provided by these instruments are local and are usually affected by the so-called “mountain breeze effect” (MBE) (e.g. Cuevas et al., 1992; Reimann et al., 2010). Therefore, these measurements are not representative of the background conditions.

Here, we present results of data that are representative of the FT using the MAX-DOAS remote sensing technique at the Izaña station, a high-mountain observatory located in the Canary Islands in the Subtropical Atlantic Ocean. Unlike in situ measurements, MAX-DOAS integrates optical paths over a few tens of kilometres. As a result, inhomogeneities are averaged in the path, which in turn minimises the impact of the small-scale upwelling of air masses from the BL. Measurements can be recorded continuously as long as diffuse sunlight is available; the best results of the proposed method are provided at solar zenith angles (SZAs) up to 70°, as discussed in Sect. 3.

Using the MAX-DOAS technique, the slant column density (SCD) of a given gas is measured, i.e. the amount of gas along a given line of sight. To remove the dependence of the observation angle from the data, the SCD is converted to the vertical column density (VCD) using a conversion factor known as the air mass factor (AMF). The AMF, which can be calculated from radiative transfer models (RTMs), is closely related to the effective optical path (EOP), which is a single path that is equivalent to all ray paths that contribute to the measurement. The EOP is needed to obtain the density at a given altitude; it can be calculated using various approaches. The simplest one was introduced by Hönninger (Hönninger and Platt, 2002; Hönninger et al., 2004). This approach only considers single scattering and does not take into account surface albedo or aerosol extinction. More realistic methods consist of using RTMs together with the optimal estimation method (OEM) (Rodgers, 2000; Steck, 2002). These methods are similar and are computationally very costly. Later, Sinreich et al. (2013) presented a new method to obtain concentrations of trace gases in the BL. The method can be applied when light paths for the lowest elevation angles do not differ significantly, which is a consequence of the high aerosol loading in the atmosphere. The optical path is then obtained from the differential SCD (DSCD) of the O₃ (Wagner et al., 2004; Sinreich et al., 2005; Frieß et al., 2006; Clémer et al., 2010), which is the difference between the SCD at low instrumental elevation angles (IEAs) and the zenith SCD. However, if the aerosol load is too high, scattering at all elevations (high and low) occurs close to the detector, which leads to a poor performance of the model. Sinreich et al. (2010) and Mahajan et al. (2012) use a similar approximation for their studies in the marine boundary layer (MBL).

In general, the different approximations used to obtain tracer concentrations from MAX-DOAS measurements require multiple scattering mode in the radiative transfer calculations due to the (a) high air density, (b) large concentrations of pollutants and often water vapour, and (c) high aerosol loading. Here, we propose a modified version of the MAX-DOAS geometric approximation (Hönninger et al., 2004) to estimate FT tracer mixing ratios at very low concentrations. In particular, at the Izaña observatory, where the air density is low and the atmosphere is extremely clear, single scattering can be used. Aerosols are not considered here, thus assumptions on the aerosol vertical distribution and properties are not necessary, and therefore uncertainties related to the aerosols are avoided. This new method, which we call the “modified geometrical approach”, takes advantage of the possibility of horizontal measurements to obtain gas mixing ratios, which provides a straightforward inversion. Here, the horizontal optical path is considered to be the difference between the optical paths of the zenith and horizon views; it is calculated using two procedures. The first procedure, called O₃-MGA (modified geometrical approach), uses the O₂ dimer (hereafter referred to as O₂) measurements at the level of the station. The second procedure, called RTM-MGA, uses a RTM to obtain the required paths. This method is valid for a single-scattering atmosphere and for scattering that occurs very close to the detector. Similar assumptions have been considered in some airborne MAX-DOAS studies (Melamed et al., 2003; Baidar et al., 2013) for nadir observations. In Baidar et al. (2013), all the photons are assumed to be scattered only once (single scattering) very close to the ground or reflected from the surface before they reach the airborne detector. In that study, the validity of these assumptions was confirmed. Here, the MGA is applied to estimate the O₃ and NO₂ concentrations at the Izaña observatory, where the required conditions are satisfied. The results are then compared with in situ measurements of both species. Because in situ measurements are based on different techniques/principles than DOAS (UV absorption cell for O₃ and chemiluminescence for NO₂), they are an independent and excellent set of measurements to validate the consistency of our method. The MGA can also be extended to other species
with structured absorption spectra in the near-UV and visible ranges, such as H$_2$O, HCHO, CHOCHO and IO.

This paper is organised as follows: in Sect. 2, the station and meteorology of the area are presented. In Sect. 3, the methodology is described and the assumptions and limitations are discussed. Sections 4 and 5 address the instrumentation and data. The results and discussion are provided in Sect. 6. Finally, the main conclusions are summarised in Sect. 7.

2 Izaña station

The Izaña Atmospheric observatory (28°18′ N, 16°29′ W) is located in Tenerife (Canary Islands) at an altitude of 2373 m a.s.l. This observatory is part of the Global Atmospheric Watch (GAW) programme and is managed by the Centro de Investigación Atmosférica de Izaña (CIAI) of the Agencia Estatal de Meteorología (AEMET, Spain). In the Canary Islands, the influence of the descending branch of the Hadley cell creates a high stability regime in the FT, whereas quasi-permanent NNE trade winds blow in the MBL (Palmén and Newton, 1969). The FT and MBL are well differentiated by a quasi-permanent strong temperature inversion that ranges between 800 and 1500 m a.s.l. (Font, 1956; Milford et al., 2008). The temperature inversion defines an upper limit to the MBL, which prevents pollution of the lower levels from reaching the FT. Therefore, the Izaña station is located well above the MBL and, during the night, is representative of the FT conditions. Below the station, at the base of the inversion, a layer of clouds is present on most days at the north side of the island. The clouds form as a result of the condensation of the humid air masses that are forced to ascend by the trade winds that arrive at the north face of the island. The Canary Archipelago is under the influence of a subtropical high-pressure system, which provides many days with a clear sky above the cloud layer.

The strong temperature inversion defines a discontinuity in the vertical distribution of gases, such as H$_2$O, NO$_2$ and O$_3$. During the night, subsiding air increases the O$_3$ and reduces the H$_2$O and NO$_2$; the opposite is true during the day. In particular, humid MBL air masses that move upslope due to the diurnal mountain breeze can reach the observatory; thus, typically, H$_2$O increases by 60 % and the NO$_2$ increases a few hundred ppt at the level of the station. In contrast, the O$_3$ suffers a decrease of 6.5 % (Cuevas et al., 1992; PuenteDura et al., 2012). These variations in concentrations can be used as signatures of the origin of the air masses over the station. As for gases, the strong temperature inversion prevents sea salt and other aerosols below from reaching the observatory. Two typical situations occur at Izaña: very clean days with an aerosol optical depth (AOD) below 0.05 at 500 nm or Saharan dust days in which the AOD can reach values larger than 1 when Saharan dust transported from North Africa arrives at the island.

**Figure 1.** Schematic representation of a simplified optical path of scattered solar radiation for real (blue) and MGA-approximated (black) geometries.

3 Modified geometrical approach (MGA)

3.1 Methodology

The MGA is an extremely simple concept. It assumes a single-scattering geometry and a scattering point altitude close to the instrument. Under these simplified conditions, the column of the gas under consideration in the horizontal path (d in Fig. 1) is the differential column between zenith (IEA = 90°) and horizontal (IEA = 0°) measurements simultaneously obtained, because the slant paths are identical in both cases and cancel out. Alternatively, the zenith measurement can be used as a reference to evaluate the horizontal measurement; then, the outcome of the evaluation is directly the gas in the horizontal path. Therefore, the concentration of the gas x at the level of the station, $c_{x,St}$, can be obtained simply from the difference of both slant columns divided by the length of the horizontal path (d):

$$c_{x,St} = \frac{SCD_{horiz,x} - SCD_{vert,x}}{d},$$

where $SCD_{horiz,x}$ and $SCD_{vert,x}$ are the slant columns in the horizontal and vertical views, respectively. If measurements are not simultaneously taken, then $SCD_{vert,x}$ has to be multiplied by a correction factor $f$ to reference them to the same SZA, which is defined as

$$f = AMF(szax1)/AMF(szax2),$$

where AMF is a factor of the zenith geometry (IEA = 90°) and szax1 and szax2 are the solar zenith angles at the times of the horizontal (IEA = 0°) and vertical (IEA = 90°) measurements.

Within the MGA, two different methods have been considered to estimate d: O$_4$-MGA and RTM-MGA. The first method obtains d from the SCD of O$_4$ (SCD$_{O_4}$), which is
measured simultaneously with the gases in which the concentration is to be retrieved and with the monthly mean oxygen concentration \((c_{O_2})\) at the level of the station. An advantage of the method is that no RTMs are needed:

\[
d = \frac{SCD_{\text{horiz}, O_4} - SCD_{\text{vert}, O_4}}{(c_{O_2})^2}.
\]

In the second method (RTM-MGA), the horizontal paths for each geometry \((d')\) are obtained from the radiative transfer model (at sza\(_1\)). Then, the gas concentration can be estimated in a similar manner as in Eq. (1):

\[
c_{x, \text{St}} = \frac{SCD_{\text{horiz}, x} - SCD_{\text{vert}, x}}{d'}.
\]

The radiative transfer model used in this paper does not directly supply the optical path. In this case, the following equation is used for the calculation of \(d'\):

\[
d' = \frac{SCD_{\text{horiz}, x} - SCD_{\text{vert}, x}}{c_x}.
\]

where \(x\) refers to O\(_3\) or NO\(_2\). In this case, the concentration of these species, \(c_x\), at the level of the station is obtained from the Air Force Geophysics Laboratory (AFGL) standard atmosphere for tropical latitudes (Anderson, 1986).

The first assumption required for the validity of the MGA is a single scattering atmosphere. To estimate the adequacy of this approximation for the conditions given in this research, the horizontal optical paths as a function of the SZA for single and multiple scattering are calculated for the Izaña observatory conditions. The SDISORT radiative transfer model (Dahlback and Stamnes, 1991) is used to obtain these horizontal paths from O\(_4\) SCDs. The inputs of these calculations are summarised in Table 2. The results are shown in Fig. 2. The differences for the SZAs that are lower than or equal to 70° are as high as 5 %. These results confirm that in the pristine Izaña atmosphere, single scattering can be used.

The MGA is also based on the assumption that when evaluating a spectrum in the horizontal direction (elevation = 0°) using the zenith spectrum as a reference, the vertical components cancel out. To prove that this is true for the conditions of this research, we use the SDISORT radiative transfer model (Dahlback and Stamnes, 1991) to calculate the single-scattering O\(_3\) differential box-AMFs defined as

\[
\text{DiffAMF} = [\text{AMF}(\alpha) - \text{AMF}(90)] \Delta z,
\]

where \(\alpha\) is a given elevation angle, and \(\Delta z\) is the thickness of the layers of constant concentration considered in the radiative transfer model (100 m in this case). This magnitude provides the difference of the optical paths for both elevation angles (0 and 90°) for each of these layers. The results are shown in Fig. 3 for an SZA = 40°. For an elevation angle = 0°, the differential box-AMF is the maximum at the level of the station but becomes zero near an altitude of 3 km and higher, which is only 0.6 km above the station. Similar results are observed for other SZAs, except for twilight (SZA > 70°), for which the method is not applicable (shown later). This result proves that the slant contribution (before the scattering point) of the paths at IEAs of 0 and 90° cancel out and that only the horizontal path remains.

### 3.2 Error estimation

The error in the concentration using Eq. (1) depends on the measurements (SCD of the considered gas) and the path errors. The average measurement relative error for typical values of NO\(_2\) and O\(_3\) and their corresponding molecular errors – as defined in Platt and Stutz, 2008 (using mean values during 2011 to 2013 at Izaña) – are 23 and 6 %, respectively. However, the path error depends on the O\(_4\) SCD fitting error (approximately 0.1 % and negligible), the air density error
(lower than 2% at Izaña in summer) and the cross-section error. Major uncertainties stem from a potential bias of the O₃ cross-sections (Wagner et al., 2002, 2009; Clemer et al., 2010); nevertheless, they are under debate. If only the air density error is considered, the path error following Eq. (3) is approximately 4%. Taking into account all these values, the typical relative errors of the NO₂ and O₃ concentrations are 23.3 and 7.4%, respectively.

The estimation of the error of the concentrations obtained from this method must also be considered. Actually, the scattering occurs at a certain height h above the station. The blue line in Fig. 1 represents the EOP, which is defined as the geometrical ray equivalent to the zenith integrated flux at the station. Accordingly, h is the “effective” scattering height. In this more realistic geometry, the difference between the horizontal and vertical measurements is

\[
\text{SCD}_{\text{horiz}} - \text{SCD}_{\text{vert}} = \text{SCD}_{\Delta}
\]

\[
= (S_1 + S_2 + D) - (S_1 + H)
\]

\[
= D + H \left( \text{AMF}_{\text{sza}} - 1 \right),
\]

(7)

where \( S_1, S_2, D \) and \( H \) are gas partial columns along the considered paths \( s_1, s_2, d \) and \( h \), respectively, and AMF(sza) = \( S_2/H \). If

\[
g = \text{AMF}_{\text{sza}} - 1,
\]

(8)

then

\[
D = \text{SCD}_{\Delta} - g H.
\]

(9)

The second term in the right-hand side of the equation is the overestimation factor. Both parameters, \( g \) and \( H \), are SZA dependent.

The relative error can be given as

\[
\varepsilon = \frac{g H}{d \bar{c}_x} = \frac{g h \tilde{c}_x}{d \bar{c}_x} = \frac{g h R}{d},
\]

(10)

where \( \tilde{c}_x \) is the mean concentration in the layer between the station and the scattering altitude, which can be written as a fraction \( R \) of the concentration at the level of the station, \( c_{x, \text{st}} \).

To estimate the relative error versus the SZA, the scattering height \( h \) must be calculated. Let \( I(z) \) be the contribution of the flux scattered at level \( z \) to the total flux at the surface. We define a weighted intensity (WI) as

\[
\text{WI}(z) = \frac{I(z)}{\int I(z) \, dz}.
\]

(11)

For a given SZA, WI represents a normalised contribution of the ray \( I \) scattered at level \( z \) to the total flux observed at the station (Solomon et al., 1987), and can be calculated with RTMs. Then, \( h \) is

\[
h = \frac{\sum_{\text{TOA}} \left( \text{WI}(z) \cdot z \right)}{\sum_{\text{station}} \text{WI}(z)}.
\]

(12)

| SZA  | Height (km) |
|------|-------------|
| 0°   | 0.0          |
| 10°  | 0.0          |
| 20°  | 0.0          |
| 30°  | 0.0          |
| 40°  | 0.0          |
| 50°  | 0.0          |
| 60°  | 0.0          |
| 70°  | 0.0          |
| 80°  | 0.0          |

Figure 4 shows that \( h \) is nearly constant up to an SZA of 70–75°. However, \( g \) increases exponentially towards large SZAs.

If the length of the horizontal path \( d \) is obtained with the help of the O₄, which is retrieved simultaneously with the NO₂ and O₃ columns (Eq. 3), then the O₄ path is also overestimated. The actual concentration of the gas \( x \) at the level of the station is

\[
c_{x, \text{st}} = \frac{\text{SCD}_{\Delta, x} - g H}{d} = (c_{O_2, \text{st}})^2 \frac{\text{SCD}_{\Delta, x} - g h \tilde{c}_x}{\text{SCD}_{\Delta, O_4} - g' h \tilde{c}_O_4}.
\]

(13)

The second terms above and below the fraction bar are the correction factors. \( \tilde{c}_x \) and \( \tilde{c}_O_4 \) are the mean concentrations in the layer between the station and the scattering altitude, which can be written as fractions, \( R \) and \( R' \), of the concentrations at the level of the station:

\[
c_{x, \text{st}} = (c_{O_2, \text{st}})^2 \frac{\text{SCD}_{\Delta, x} - g h R c_{x, \text{st}}}{\text{SCD}_{\Delta, O_4} - g' h R' (c_{O_2, \text{st}})^2},
\]

(14)

which, with reassigning, yields

\[
c_{x, \text{st}} = \frac{\text{SCD}_{\Delta, x}}{(c_{O_2, \text{st}})^2 + h(R - R')},
\]

(15)

The error in the path, the term \( h(R - R') \), depends on the vertical distribution of the tracer under consideration and the differences in the AMF between the tracer and the O₄. If the vertical distribution does not vary much with height, we can estimate the error as differences in the AMF. Figure 5 shows the relative error in O₃ and NO₂ assuming a horizontal path of 60 km (approx. AOD⁵₀₀ nm = 0.02) and a constant vertical volume mixing ratio (VMR). As observed, the best results of the MGA are obtained for SZAs up to 70°.

Note that the observed errors do not increase if a moderate amount of aerosols are present (AOD⁵₀₀ nm = 0.1) because both the numerator and denominator in Eq. (15) are reduced by approximately the same ratio.
4 MAX-DOAS measurements and data

MAX-DOAS measurements were recorded at a terrace on the Izaña station using an INTA spectrometer, the RASAS-II, in the summer of 2011. In Table 1, the instrument settings are summarised. Further details can be found in Roscoe et al. (2010) and Puentedura et al. (2012). Note that the field of view of the light collector was narrowed (with respect to previous studies) to 1°, which reduced the uncertainties of the scanned air masses. Most of the instrument-collected radiation originated from the FT (Puentedura et al., 2012), and the largest contribution was from the station level. RASAS-II is pointed north and collects scattered radiation from several elevation angles between −1 and 90°. However, because the sun is directly overhead at noon during the subtropical summer, the zenith data were not collected at that time to avoid damage due to direct sunlight on the detector. Data at IEA = 70° are instead used in our calculations. Consequently, SCD_{vert} in the equations corresponds to α = 70° instead of 90°. The impact of this change was tested and found to be negligible.

Spectra recorded by RASAS-II were processed to obtain the SCDs of O₃, O₄ and NO₂. This retrieval was performed using a code developed at INTA (Gil et al., 2008) based on the standard DOAS technique (Platt and Stutz, 2008). In the present analysis, a spectral range from 430 to 500 nm was chosen. This range differs from that suggested by the Network for the Detection of Atmospheric Composition Change (NDACC) for O₃ [450–550 nm]. An extension to 430 nm was chosen to perform a single NO₂ and O₃ analysis that covers the large O₄ 477 nm absorption. The longest wavelengths were truncated to avoid the strong H₂O absorption band at 500–510 nm. H₂O bands cannot be accurately fitted because the effective absorption cross sections are dependent on the actual H₂O concentration due to saturation (Clémer et al., 2010). Although the truncation of the spectral upper range at 500 nm improves the fitting, an inspection of the remaining residuals clearly shows systematic structures in the flanks of the water cross-sections’ weak bands centred at ∼445 and ∼487 nm and the strong 477 nm O₄ band. The structures were removed using a pseudo-cross-section generated as the residual mean of 3 days of clear data (approx. 1000 spectra). Careful testing was performed to ensure that no interference occurred. The results had almost no impact on the retrieved amounts, and the stability of the IEA cycles (differences between the contiguous cycles) improved.

Although the differential absorption signal between 70° and the horizon was small in the unpolluted FT, the measurements were of good quality. Following Platt and Stutz (2008), the typical instrumental detection limits (DL) for the measurements used in the present study were 4 × 10¹³ and 8 × 10¹⁶ molecules cm⁻² for NO₂ and O₃, respectively, for the typical root mean square error of the residuals of a 2.5 × 10⁻⁴ differential optical depth (DOD). Assuming no error occurred in the optical path estimation, this DL was lower than 1 ppt for NO₂ and 1 ppb for O₃. As mentioned previously, the contribution of the fitting error of O₃ was negligible (0.1 %). Major uncertainties stem from the potential bias.
in the O\textsubscript{4} cross-sections (Wagner et al., 2002, 2009; Clémer et al., 2010) used to obtain the optical path. An example of the DOAS fitting is shown in Fig. 6, and the standard diurnal evolution of DSCDs of O\textsubscript{3}, NO\textsubscript{2}, O\textsubscript{4} and H\textsubscript{2}O on a clear day for all IEAs is shown in Fig. 7. Note that NO\textsubscript{2} and O\textsubscript{3} DSCD plots have a u-shape due to their high concentration in the stratosphere. This signature is not observed in the DSCDs of O\textsubscript{4} because most of the oxygen resides in the troposphere.

5 In situ measurements

The performances of O\textsubscript{4}-MGA and RTM-MGA were tested by comparing their estimated concentrations of O\textsubscript{3} and NO\textsubscript{2} with the in situ measurements available at the Izaña observatory. The NO\textsubscript{2} and O\textsubscript{3} in situ concentrations were measured with a chemiluminescence NO-NO\textsubscript{2}-NO\textsubscript{X} analyser (Model 42C-TL, Thermo Electron Corporation) and two UV photometric O\textsubscript{3} analysers running in parallel (Model 49C, Thermo Electron Corporation), respectively. Detection limits of these instruments are 50 ppt and 1 ppb, respectively. The sampled air masses were captured by two inlet systems on the top of the observation tower, which is located 4 m above the terrace. One of the inlets is used for the O\textsubscript{3} determination, and the second inlet is used for measuring the remaining tracers. A laminar vertical flux manifold ensures that the residence times of the air masses along the sample are smaller than 10 s. The quality control for ozone (Cuevas et al., 2013) includes a 15 min daily check of zeros, calibration of the O\textsubscript{3} analyser with a primary standard model 49C-PS, and quality audits periodically performed by the World Calibration Centre (WCC-EMPA). The calibration of the NO\textsubscript{2} analyser (González, 2012) is performed using a multi-gas calibrator and certified gas bottles provided by Air Liquide Inc. The Environmental Protection Agency of the United States (EPA) and the European Union have declared these measurement techniques as reference methodologies. Furthermore, these techniques fulfil the Global Atmospheric Watch Programme requirements.

The main drawback of the measurements provided by these instruments is that they are usually affected by the radiatively induced local circulation. During the day, the upslope wind transports tracers from the lower levels (MBE, e.g. Cuevas et al., 1992; Reidmiller et al., 2012), whereas at night, a quasi-permanent katabatic NW flow is normally present at the level of the station, except during Saharan air mass intrusions (Cuevas et al., 2013). As a result, only night-time in situ measurements are representative of FT conditions. The diurnal intensity of the MBE is modulated by insolation and the synoptic wind. The stronger the synoptic wind, the smaller the influence of the MBE (Cuevas et al., 2013).

Ozone sounding data of the NDACC programme are also available. Ozonesondes are launched on a weekly basis from the Orotava station (28°25′N, 16°18′W) in Puerto de la...
Cruz (Tenerife) at a horizontal distance of approximately 13 km from Izaña. Ozone soundings provide O$_3$ concentrations and meteorological data (pressure, temperature and humidity) from sea level to 30–35 km altitudes with a vertical resolution of 100–150 m (Cuevas et al., 1994). O$_3$ is measured by the standard ECC method (Thompson et al., 2011; Smit, 2011).

6 Results and discussion

The O$_3$ and NO$_2$ concentrations at the Izaña station were estimated using both O$_4$-MGA and RTM-MGA. To ensure that the assumptions of the approximation are met (single-scattering atmosphere and scattering close to the detector), only measurements at SZAs lower than 70° were used. Measurements taken on clear days (AOD below 0.05) were considered here. Figure 8 shows the AERONET level 2 (cloud-screened and calibrated data) AOD at 500 nm obtained by the CIMEL sun photometer over the station during July and the beginning of August 2011. Five well-identified Saharan events occurred during this period. The rest of the days exhibited clear skies with typical AODs below 0.05. In this period, we selected a set of 9 summer days (days 200–208 of 2011) in which the first 8 correspond to clear sky days. On the last day, only the first hours were considered because Saharan dust arrived at the station in the early afternoon.

O$_3$, NO$_2$ and O$_4$ SCDs were simultaneously measured by the MAX-DOAS spectrometer. For the O$_4$-MGA, the path was directly obtained by Eq. (3). Because the pressure variation from one day to another is only 2% in summer, the air...
density radiosonde monthly climatology was used to obtain the O\textsubscript{4} concentration at the level of the station.

Optical paths were also estimated from the radiances obtained by the pseudo-spherical discrete ordinate solver SDISORT (Dahlback and Stamnes, 1991) that is included in the software package libRadtran version 1.7 (Mayer and Kylling, 2005) for the RTM-MGA. This solver considers direct radiation in a spherical Earth framework and the scattering of light in a plane-parallel configuration. Vertical profiles of the gas densities were taken from the AFGL standard atmosphere for tropical latitudes (Anderson, 1986). Due to the low AOD found at the station for the selected days, no aerosols were considered in the simulations. Furthermore, to best reproduce the conditions of the site, a homogeneous cloud located between 600 and 1100 m a.s.l. was included in the model. The altitude values of the cloud top and base were estimated from the European Centre for Medium-Range Weather Forecasts (ECMWF) using summer conditions from 2008 to 2011 for 24 h forecasts computed at intervals of 6 h. These altitude values are consistent with the radiosoundings and present a small day-to-day variability during summer (J. J. Bustos, personal communication, 2010). Additionally, we used an optical depth of the cloud equal to 5 (at 500 nm), as catalogued by the International Satellite Cloud Climatology Project (IS-CCP), and a standard marine surface albedo of 0.07. The instrumental azimuth is set to zero, and the solar azimuth dependence is considered in all of our calculations. The RTM calculations were performed at 440 nm for O\textsubscript{3} and NO\textsubscript{2}. Under these conditions (summarised in Table 2), the values of the difference between the optical paths for the zenith and horizon views (\textit{d} in the equations) were obtained for SZAs from 0 to 70° in intervals of 1°. To obtain intermediate values, linear interpolation was used.

### 6.1 Ozone concentration

In Fig. 9, the O\textsubscript{3} mixing ratios obtained using O\textsubscript{4}-MGA (blue) and single-scattering RTM-MGA (red) are compared to the O\textsubscript{3} in situ measurements (black). The concentrations based on O\textsubscript{4}-MGA yield diurnal evolutions similar to those obtained using RTM-MGA. The MGA captures the day-to-day variability well. However, the mean difference between the in situ and the MAX-DOAS values for the selected period was 28 %. In general, larger values are observed by MAX-DOAS. The possible causes of this discrepancy are described below.

The O\textsubscript{4}-MGA relies on the accuracy of the O\textsubscript{4} absorption cross-sections that are known to have large uncertainties (Wagner et al., 2002). Even at room temperature, the O\textsubscript{4} cross-sections reported in the literature (Greenblatt et al., 1990; Hermans et al., 1999) seem to provide an overestimation (up to 30 %) of the DSCD (Wagner et al., 2002, 2009; Clémer et al., 2010), but the magnitude of the correction to reconcile the models and measurements seems to be dependent on a number of factors. A case study was performed over Izaña on day number 204 at the O\textsubscript{4} 477 nm band; a good correlation (\textit{R} \textsuperscript{2} = 0.96) was identified between the measurements and the RTM, which provides confidence in the magnitude of the retrieved optical paths. A similar result has been obtained recently by Baidar et al. (2013) using high-altitude aircraft measurements, which suggests that pressure, temperature or both play a role in the magnitude of the observed discrepancy. At our station, smaller O\textsubscript{4} cross-sections should be required to shift the measured O\textsubscript{3} concentrations towards the in situ values.

The differences between the MGA and in situ measurements may also be due to the different air masses scanned by each instrument. MAX-DOAS scans a FT area over the ocean far from the surface along a 60–70 km path. Therefore, the MAX-DOAS spectrometer provides the O\textsubscript{3} background in the FT. During summer, a subtropical high-pressure system induces subsidence that supplies high O\textsubscript{3} air masses from above, whereas at the observatory the effect is the opposite. A radiatively induced upslope breeze transports O\textsubscript{3}-rich air from lower altitudes, which reduces the O\textsubscript{3} concentration. A similar result has been obtained recently by Baidar et al. (2013) using high-altitude aircraft measurements, which suggests that pressure, temperature or both play a role in the magnitude of the observed discrepancy. At our station, smaller O\textsubscript{4} cross-sections should be required to shift the measured O\textsubscript{3} concentrations towards the in situ values.

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### Table 2. Input parameters used in libRadtran for the optical path calculations.

| Parameter               | LibRadtran inputs |
|-------------------------|--------------------|
| Atmosphere model        | Standard atmosphere for tropical latitudes: afgl |
| Aerosols                | No aerosols       |
| Cloud                   | Optical depth 5, 0.6–1.1 km |
| Albedo                  | 0.07               |
| Instrument azimuth      | Fixed to north    |
| Solar azimuth           | Variable          |
| Solar zenith angle      | From 0 to 70° in 1° intervals |
| Elevation angle         | 0.1 and 70°       |
| Wavelength              | 440 nm (O\textsubscript{3} and NO\textsubscript{2}), 477 nm (O\textsubscript{4}) |
| Scattering              | Single scattering |
| Cross-section (O\textsubscript{3}) | Bogumil et al. (2001) (223 K) |
| Cross-section (NO\textsubscript{2}) | Vandaele et al. (1998) (220 K) |
| Cross-section (O\textsubscript{4}) | Greenblatt et al. (1990) (296 K) |
| Solver                  | SDISORT            |
The O$_3$ data at the level of the observatory measured by the ozone sounding launched on day 208 also supports this possibility (i.e. the different scanned air masses as the cause of the difference between MGA and in situ O$_3$ mixing ratios), see Fig. 9. Considering that the ECC O$_3$ sonde accuracy in the troposphere is 8–10% (Smit, 2011) (~8 ppbv on day 208 for an 80 ppbv mixing ratio at the level of the station), the O$_3$ value from the ozone-sonde at the Izaña level is in rather good agreement with the estimated MGA values. Note that during the study period (days 200–208), the in situ measurements end around noon for technical reasons. As a result, the MAX-DOAS and in situ measurements overlap only during some daylight hours. This could also be a cause of the observed differences between the MGA and in situ O$_3$ concentrations. To test this hypothesis, the mean diurnal O$_3$ volume mixing ratios for the RTM-MGA and in situ measurements were calculated. The daily averages for each instrument were calculated from two sets of data: all available data for each day and data for the periods of the day when simultaneous measurements were available. The results show that, when considering only simultaneous data, the differences between the MGA and in situ approach are not reduced, which eliminates the in situ noon blackouts as the cause of the discrepancy.

On some of the analysed days, the diurnal evolution of the O$_3$ concentration shows a decrease during the day that seems too large to be genuine. This is not observed for O$_4$ or NO$_2$, which suggests that the variation is not optical path related (i.e. azimuth changes). New experiments are in progress to identify the causes of this effect.

### 6.2 NO$_2$ concentration

The same procedure that was applied to O$_3$, was also used to obtain the NO$_2$ concentrations, but hourly profiles were calculated to account for photochemical changes. This allowed us to test the sensitivity of the retrieved concentration to the profile. The hourly profiles were obtained using a photochemical box model (Denis et al., 2005) derived from the SLIMCAT 3-D chemical transport model (Chipperfield, 2006).

In Fig. 10, the NO$_2$ results of RTM-MGA (red) are shown together with the in situ data (black). The O$_4$-MGA results are shown as blue diamonds. The NO$_2$ VRM are at the 20–30 pptv level. For such low NO$_2$ concentrations, no tools are available to decide which of the two retrievals is the most accurate. Nevertheless, the in situ NO$_2$ measurements are shown for reference because, during the night-time, the NO$_2$ concentrations are below the DL given by the manufacturer (50 ppt for the 300 s average).

Contrary to the observations of the O$_3$, the in situ NO$_2$ measurements increase to very high concentrations during the hours around noon. This increase is the result of the MBE. The mountain breeze brings anthropogenic NO$_2$, which originated over the populated areas of the coast, to the station. Thus, the mixing ratios increase to hundreds of ppt.
Figure 10. Surface NO$_2$ mixing ratios at the Izaña observatory for days 200 to 207. The black crosses correspond to the in situ measurements. The blue diamonds correspond to the NO$_2$ mixing ratios obtained with the O$_4$-MGA. The red solid circles correspond to the NO$_2$ mixing ratios obtained with the RTM-MGA.

around noon (Volz-Thomas et al., 1993; Puentedura et al., 2012). This local effect, which renders the in situ data non-representative of the FT during the daytime, is no more than a few hundred metres above the surface. The relative contribution of this narrow BML layer to the MAX-DOAS path is very low, and it is smoothed out in the long paths that were sampled. At night, a well-developed katabatic regime transports air masses from higher altitudes, which ensures that the measured NO$_2$ is representative of the FT conditions.

The NO$_2$ concentrations obtained from the MGA, as presented in this study, were also compared with previous FT NO$_2$ measurements obtained from research aircraft flights. Bucsela et al. (2008) measured NO$_2$ over the eastern Atlantic coast of North America during the Intercontinental Chemical Transport Experiment – North America, Phase A (INTEX-A), and the International Consortium for Atmospheric Research and Transformation (ICARTT) aircraft campaigns using the research DC-8 of NASA. These measurements were obtained using the UC Berkeley laser-induced fluorescence instrument (TD-LIF) whose DL is 4 ppt. These profiles reveal NO$_2$ concentration values that are similar to those found in this study (approximately 20 ppt) for the same altitude and the same time of the year (Martin et al., 2006; Bucsela et al., 2008).

7 Summary and conclusions

A new method, the MGA, was developed to obtain mixing ratios of trace gases in the FT. The MGA uses horizontal and near-zenith geometries to estimate a station level differential path. Two independent methods were used for the optical path calculation: the first method, O$_4$-MGA, uses MAX-DOAS SCDs of O$_4$; the second method, RTM-MGA, obtains optical paths from AMFs obtained optical paths from a RTM. The methodology is useful at high mountain observatories where horizontal measurements are possible and aerosol loading is low. The validity of the assumptions considered in this approximation and the applicability of this method have been confirmed for SZAs lower than 70$^\circ$.

Both methods, O$_4$-MGA and RTM-MGA, were applied to MAX-DOAS measurements at the Izaña observatory to estimate FT O$_3$ and NO$_2$ mixing ratios. By comparing these results with in situ measurements, it is shown that under low aerosol loading, the mixing ratios of O$_3$ and NO$_2$ can be retrieved. Day-to-day O$_3$ variations are well captured by the MGA, although mean differences of 28% between the MGA results and in situ measurements were observed. The different air masses that were scanned by each instrument are possible causes of these discrepancies. The NO$_2$ concentrations obtained with the MGA are within the range of 20–40 ppt, which is below the DL of the in situ chemiluminescent analyser but in good agreement with other subtropical FT measurements recorded by aircraft (Martin et al., 2006; Bucsela et al., 2008).

The main advantage of the MGA is that it provides mixing ratios that are only slightly affected by the MBE. Because air masses sampled by a MAX-DOAS spectrometer extend a few tens of kilometres, inhomogeneities are smoothed along the path. Therefore, unlike in situ data, MAX-DOAS
measurements are representative of the FT. Furthermore, the 
MGA provides a simple way to estimate O$_3$ mixing ratios 
when no in situ instruments are available. In fact, for NO$_2$, 
the method provides mixing ratios of a few tens of pptv, 
which are below the DL of the standard instrumentation 
based on chemiluminiscence.

Acknowledgements. Laura Gomez thanks the MINECO (Spanish 
Economy Ministry) for providing the Jaan de la Cierva grant. 
We acknowledge the support of AMISOC (Atmospheric Minor 
Species relevant to the Ozone Chemistry at both sides of the 
Subtropical jet, contract number CGL2011-24891), and NORS 
(Demonstration Network of Ground-Based Remote Sensing 
Observations in support of the GMES Atmospheric Service) 
integrated project under the 7th Framework Programme (contract 
number FP7-SPACE-2011-284421). The authors thank Ramón 
Ramos and the CIAI staff for their support at the Izaña observatory 
and the Sieltec team for the instrument maintenance. We thank 
Philippe Goloub for his effort in establishing and maintaining the 
AERONET Izaña site.

Edited by: M. Van Roozendael

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