GAMMA-RAY OBSERVATIONAL PROPERTIES OF TeV-DETECTED BLAZARS

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ABSTRACT

The synergy between the Fermi-LAT and ground-based Cherenkov telescope arrays gives us the opportunity for the first time to characterize the high-energy emission from blazars over 5 decades in energy, from 100 MeV to 10 TeV. In this study, we perform a Fermi-LAT spectral analysis for TeV-detected blazars and combine it with archival TeV data. We examine the observational properties in the γ-ray band of our sample of TeV-detected blazars and compare the results with X-ray and GeV-selected populations. The spectral energy distributions (SEDs) that result from combining Fermi-LAT and ground-based spectra are studied in detail. Simple parameterizations such as a power-law function do not always reproduce the high-energy SEDs, where spectral features that could indicate intrinsic absorption are observed.

Key words: galaxies: active – galaxies: nuclei – gamma rays: general

Online-only material: color figures

1. INTRODUCTION

Active galactic nuclei (AGNs) are extreme objects with observed luminosity outshining their host galaxy. These sources are believed to be powered by accretion onto a central supermassive black hole, commonly display relativistic jets, and exhibit non-thermal continuum emission extending from the radio band to X- and γ-rays. Blazars constitute a subclass of AGNs, with jet axes oriented close to the observer’s line of sight. Relativistic beaming gives rise to distinctive observational features in blazars, such as strongly anisotropic radiation, superluminal motion, high polarization, and rapid variability (Urry & Padovani 1995). Blazars are divided into two subclasses, flat spectrum radio quasars (FSRQs) and BL Lacertae objects (BL Lacs). FSRQs are observationally characterized by broad spectral lines in the optical band, which are weak or not present in BL Lacs. The spectral energy distribution (SED) of blazars exhibits a two-component structure, with a low-energy component peaking between infrared (IR) and X-ray energies, and a high-energy one between X- and γ-rays. The low-energy component is believed to be dominated by synchrotron emission from relativistic electrons in the jet (Kembhavi & Narlikar 1999). The peak frequency of the synchrotron component of the SED ($\nu_{\text{syn}}$) is used to subclassify BL Lacs into low (LBLs; $\nu_{\text{syn}} < 10^{14}$ Hz), intermediate (IBL; $\nu_{\text{syn}} \sim 10^{14}$-$10^{15}$ Hz), and high-frequency-peaked BL Lacs (HBL; $\nu_{\text{syn}} > 10^{15}$ Hz).

The high-energy component of the blazar SED has been historically less studied, due to the later development of hard X-ray and γ-ray detectors compared to those of longer frequency bands. The synchrotron self-Compton (SSC) model is the simplest scenario that explains the high-energy emission of blazars, by inverse-Compton (IC) upscattering of soft synchrotron photons off the same electrons that have undergone synchrotron cooling (Maraschi et al. 1992). Throughout the text, we refer to the high-energy component of the blazar SED as “IC component.” An additional IC target photon field external to the jet is often invoked (see, e.g., W Comae: Acciari et al. 2009d). This mechanism is referred to as external-Compton (EC), and several possible sources for the external photon field have been set forth (Böttcher 2010). Other models suggest a significant contribution from hadronic processes to the high-energy output (Mannheim & Biermann 1992).

A good spectral characterization of the high-energy peak of the blazar spectrum (keV–TeV band) is essential to discriminate between the aforementioned models. During the EGRET era covering 1991–2001 (Thompson et al. 1993), only five blazars were known at TeV energies: Mrk 421, Mrk 501, 1ES 1959+650, PKS 2155–304, and 1ES 2344+514; thanks to the first generation of ground-based Imaging Atmospheric Cherenkov Telescopes (IACTs; Whipple, Kildea et al. 2007; HEGRA, Pühlhofer et al. 2003; Durham Mark 6, Armstrong et al. 1999; and Telescope Array, Aiso et al. 1997). Only three of these sources were detected by EGRET (Mrk 501 was only marginally detected and 1ES 2344+514 was not seen at all). By the time Fermi started operations (2008 August), the number of known TeV blazars had increased to 21 with the second generation of IACTs in operation (VERITAS, Weekes et al. 2010; MAGIC, Aleksić et al. 2012; HESS, Bernlörhr et al. 2003). This number has doubled since then, with most TeV blazars also being detected in the GeV range by Fermi-LAT.

For the first time now, good quality spectra are available both from Fermi-LAT in the high-energy (HE; 0.1 GeV < $E$ < 100 GeV) γ-ray band and IACTs in the very high energy (VHE; $E$ > 100 GeV) γ-ray band for more than two dozen sources. The combined spectral data cover up to five decades in energy, giving a detailed description of the high-energy peak of the blazar SED. Recent studies have explored this newly available data sample, focusing on the GeV properties of TeV-selected blazars (Abdo et al. 2009), or deriving jet parameters assuming leptonic emission models (Zhang et al. 2012). These studies are similar to earlier studies carried out on a limited sample of TeV-detected blazars (e.g., Wagner 2008).

In this paper we study the GeV–TeV observational properties of the high-energy emission in blazars that are detected in the TeV band. Section 2 describes the population of TeV blazars,
Figure 1. Sky map of TeV blazars in galactic coordinates, as of 2012 January, generated using TeVCat (http://tevcat.uchicago.edu/). The blue and pink shaded areas represent VERITAS/MAGIC and HESS visibilities, respectively. A total of 46 sources consisting of 33 HBLs, 4 IBLs, 4 LBLs, 3 FSRQs, and 2 sources that were formerly classified as AGNs of unknown type (UNID), namely IC 310 and VER J0521+211, are shown. VER J0521+211 is now identified as a BL Lac (Errando et al. 2011b) and recent studies suggest that the high-energy radiation from IC 310 originates from a blazar-like emission mechanism (Kadler et al. 2012).

(A color version of this figure is available in the online journal.)

giving census information, investigating luminosity, redshift, and photon index distributions among different blazar types. In Section 3, we study TeV blazars that appear in the Fermi data and outline their GeV properties with respect to the rest of the Fermi blazars. Section 4 defines our sample and focuses on general observational TeV properties of our objects. In Section 5, we give a detailed description of the Fermi analysis that we performed on our TeV blazar sample. Finally, Section 6 discusses various observational characteristics of the studied sources based on their GeV–TeV spectral shapes, such as the peak frequency of the IC component, absorption-like spectral features, and variability. Throughout the text, the symbol σ is used to designate the standard deviation, as a measure of statistical significance.

2. TeV BLAZARS

Mkn 421 was the first blazar and extragalactic object to be discovered as a VHE γ-ray emitter, detected with the Whipple telescope in 1992 (Punch et al. 1992). Since then, different candidate selection methods have been applied to radio, X-ray or HE data with the aim of finding new “TeV” blazars, i.e., detected in the VHE regime (Costamante & Ghisellini 2002; de la Calle Pérez et al. 2003; Behera & Wagner 2009), leading to the discovery of most of the known TeV blazars. To date, 44 blazars and 2 AGNs of unknown type have been detected in the VHE range,6 with a census consisting of 33 HBLs, 4 IBLs, 4 LBLs, and 3 FSRQs (see Figure 1). In this work, we have studied the blazars that have a published TeV spectrum as of 2011 February (referred to as the “sample” in the remainder of the text).

The redshift (z) of TeV blazars in our sample ranges from 0.031 (HBL Mrk 421) to 0.536 (FSRQ 3C 279), and nearly one-fourth of the population does not have a secure redshift. This lack is due to the fact that optical emission lines are typically weak or absent in BL Lac objects, rendering direct redshift measurements difficult. The majority of known-redshift TeV blazars are located at $z < 0.2$, mostly due to the absorption from the extragalactic background light (EBL). Figure 2 illustrates the redshift distribution of all TeV blazars. The TeV FSRQs are the most distant objects in the population. The farthest object is the FSRQ 3C 279, with a redshift of 0.536. The rest of the population does not have a secure redshift.

(A color version of this figure is available in the online journal.)

6 http://tevcat.uchicago.edu/
for $E > 400$ GeV, with the following formula:

$$L = F \times 4\pi D_L^2/(1 + z)^2 - \Gamma,$$

(1)

where $F$ is the energy flux for energies above $E > 400$ GeV, $D_L$ is the luminosity distance with Hubble constant $H_0 = 71$ km s$^{-1}$ Mpc$^{-1}$ and the cosmological constant $\Omega = 0.730$, $z$ is the redshift, and $\Gamma$ is the observed photon index for each blazar. Figure 3 (left) shows the luminosity versus redshift correlation for the sample. Sources at high redshifts tend to be scarce and much more luminous. The reason why we see only luminous sources at high redshifts is that the less luminous ones are too weak to be detected. On the other hand, the reason why we do not detect them in low redshifts is that we integrate over a much smaller volume and thus are less likely to see high-luminosity sources that should be scarce compared to low-luminosity ones. Note that if the luminosity is corrected for EBL absorption, which is stronger at TeV energies and high redshifts, the correlation will be steeper. Figure 3 (right) shows the TeV photon index versus luminosity correlation of the same sample. For luminosities up to $\sim 10^{45}$ erg s$^{-1}$, the photon index distribution is fairly homogenous.

3. Fermi TeV BLAZARS

The second Fermi-LAT catalog (2FGL) contains 1873 sources, among which 1062 are AGNs, with 435 BL Lacs, 370 FSRQs, and 257 AGNs of unknown class (Nolan et al. 2012). Thirty-six of these AGNs are TeV emitters. Figure 4 (left) shows the distribution of Fermi spectral indices ($\Gamma_{\text{GeV}}$) for TeV and non-TeV blazars in the 2FGL catalog. TeV-detected blazars tend to have harder GeV indices. As can be seen from Figure 4 (right), another distinguishing parameter for TeV emitters within the Fermi blazar population is the integral flux for energies above 1 GeV. It follows that an effective method for TeV-candidate selection in the HE $\gamma$-ray band is to look for bright hard spectrum sources and select the candidates based on their extrapolated fluxes at VHE energies. For all TeV blazars, $\Gamma_{\text{GeV}} < 2.3$ and for most of them $\Gamma_{\text{GeV}} < 2$, in agreement with an inverse-Compton peak frequency ($\nu_{\text{IC}}$) located in the high-energy tail of the Fermi range or beyond. Figure 5 shows scatter plots of spectral indices of Fermi-bright AGNs (Abdo et al. 2010a) in radio, optical and X-ray bands, comparing TeV and non-TeV sources. The TeV and non-TeV AGNs occupy separate regions in the parameter space, consistent with the results in Abdo et al. (2010a), considering that most TeV AGNs are HBLs.
Figure 5. Scatter plots of $\alpha_{ro}$ (5 GHz–5000 Å), $\alpha_{ox}$ (5000 Å–1 keV) and $\alpha_{x\gamma}$ (1 keV–100 MeV) spectral indices for the Fermi-LAT bright AGN sample (LBAS). Data are taken from Abdo et al. (2010a). Note the correlation between $\alpha_{ro}$ and $\alpha_{x\gamma}$. The red triangles represent TeV-detected AGNs, and asterisks non-TeV-detected ones. The TeV-emitters seem to be well isolated in $\alpha_{ro}$–$\alpha_{x\gamma}$ parameter space. (A color version of this figure is available in the online journal.)

Figure 6. TeV photon index distributions for our sample (see Table 1). For blazars that have multiple published results in Table 1, the most recent one was used. Top, middle, and bottom panels show HBLs, IBLs+LBLs, and FSRQs, respectively. HBLs tend to have harder spectra than the rest of the sample. FSRQs have the softest spectra. Note that the TeV indices are not EBL-corrected. (A color version of this figure is available in the online journal.)

4. DATA SAMPLE

Our blazar sample contains all blazars with a published VHE spectrum before 2011 February, including a total of 26 sources (see Table 1): 19 HBLs, 3 IBLs, 2 LBLs, and 2 FSRQs. TeV spectral index distributions of the whole sample are shown in Figure 6. Three of these blazars have insecure redshifts either because the spectroscopic measurements were inconclusive, or the calculations were made indirectly based on EBL absorption studies. References for the adopted redshift values in these three cases are given in Table 1. Seven of our targets were detected with EGRET and 23 of them are in the 2FGL catalog (Nolan et al. 2012). The ones that are missing in the Fermi data (1ES 0229+200, 1ES 0347-121, PKS 0548-322) are very hard spectrum sources that would be weak in the Fermi band. More than half of the sample have been detected multiple times in the VHE band. These multiple detections extending over several years and obtained mostly with different instruments suggest that spectral variability in the VHE band is a common property for VHE blazars. Even though no general pattern has been established for VHE variability, several sources have been observed to have a flux increase up to a few times their baseline emission (W Comae: Acciari et al. 2009d; Mkn 501: Albert et al. 2007e; PKS 2155-304: Aharonian et al. 2007d), occasionally accompanied by a change in spectral index (Mkn 501: Albert et al. 2007e) and minute-scale flux doubling times (Mkn 501: Albert et al. 2007e; PKS 2155-304: Aharonian et al. 2007d).

The first 27-month Fermi data and archival VHE spectra published before 2011 February were used to construct combined GeV–TeV SEDs in this study. Only in seven cases (RGB J0710+591, 1ES 1218+304, PKS 1222+21 (4C+21.35) PKS 1424+240, PKS 2155-304, and two different measurements for 3C 66A) were the VHE data found to overlap with the Fermi era. The remainder of the VHE data were taken before the Fermi mission.

All VHE spectra were corrected for the EBL absorption using the model by Domínguez et al. (2011). Other background models are also available (e.g., Finke et al. 2010). However, with a different EBL model, we do not expect any significant differences in our results up to a few TeV, given the redshift and energy range of our sample. See Section 6.2 for a more detailed discussion on the EBL correction effects on our study.

5. Fermi ANALYSIS

The fact that most of the GeV and TeV data are not contemporaneous makes it hard to interpret the combined spectra of blazars. Moreover, Fermi data represent an average state over relatively long periods, whereas the VHE spectra consist of “snapshots,” mostly taken during flares. To account for blazar variability and the non-contemporaneous nature of the data set, for bright enough sources, the Fermi data were split into “low” and “high” flux states as described below. Thus, non-contemporaneous GeV and TeV measurements were matched in a more realistic way than directly using all the time-averaged Fermi data. Table 1 summarizes the Fermi flux states and VHE spectra used for each source.
Table 1

| Name                  | SED Type | $\alpha$ | $\log_{10}(\text{flux}/1\text{ TeV})$ | Fermi Var. | Fermi State | $\Gamma_{\text{TeV}}$ | $I$ (7) | $E_{\text{in}}$ (GeV) | Reference (9) |
|-----------------------|----------|----------|---------------------------------------|------------|--------------|-----------------------|--------|----------------------|----------------|
| RGB J0152+017         | HBL      | 0.080    | ...                                   | ...        | average      | 2.95 ± 0.36         |        | 2%                   | (Aharonian et al. 2008) |
| 3C 66A*               | BRL      | 0.444*   | 15.63                                 | 171        | MAGIC        | 3.64 ± 0.39         | 6%     | 2%                   | (Aleksić et al. 2011a) |
| iES 0229+200          | HBL      | 0.140    | 19.45                                 | ...        | average      | 2.50 ± 0.19         |        | 2%                   | (Aharonian et al. 2007b) |
| iES 0347-121          | HBL      | 0.188    | 17.94                                 | ...        | average      | 3.10 ± 0.23         |        | 2%                   | (Aharonian et al. 2007a) |
| PKS 0548-322          | HBL      | 0.069    | 16.84                                 | ...        | average      | 2.86 ± 0.34         | 1%     |                      | (Aharonian et al. 2010) |
| RGB J0701+591         | HBL      | 0.125    | 21.05                                 | 6          | VERITAS      | 2.69 ± 0.26         | 3%     |                      | (Acciari et al. 2010c) |
| S5 0716+714           | LBL      | 0.300    | 14.46                                 | 266        | high         | 3.45 ± 0.54         | 9%     |                      | (Anderhub et al. 2009a) |
| iES 0806+524          | HBL      | 0.138    | 16.56                                 | 20         | average      | 3.6 ± 1.0           |        | 2%                   | (Acciari et al. 2009a) |
| iES 1011+496          | HBL      | 0.212    | 16.74                                 | 16         | high         | 4.0 ± 0.5           |        | 6%                   | (Albert et al. 2007d) |
| iES 1101-232          | HBL      | 0.186    | 16.88                                 | 1         | average      | 2.94 ± 0.20         |        | 3%                   | (Aharonian et al. 2007c) |
| Markarian 421*        | LBL      | 0.031    | 18.49                                 | 44         | medium       | 2.20 ± 0.08         | 50–200% | 200%                 | (Albert et al. 2007b) |
| Markarian 180         | HBL      | 0.046    | 18.61                                 | 10         | average      | 3.3 ± 0.7           |        | 11%                  | (Albert et al. 2006) |
| iES 1218+304*         | HBL      | 0.182    | 19.14                                 | 15         | average      | 3.08 ± 0.34         |        | 7%                   | (Acciari et al. 2009c) |
| W Comae*              | BBL      | 0.102    | 18.44                                 | 47         | high         | 3.81 ± 0.35         |        | 9%                   | (Acciari et al. 2008) |
| PKS 1222+21 (4C+21.35)| FSRQ     | 0.432    | 13.27                                 | 101        | MAGIC        | 3.75 ± 0.27         | 100%   | 4%                   | (Aleksić et al. 2011b) |
| 3C 279                | FSRQ     | 0.536    | 12.67                                 | 898        | high         | 4.1 ± 0.7           |        | 15%                  | (Albert et al. 2008) |
| PKS 1424+240          | BBL      | 0.260*   | 15.7                                  | 26         | VERITAS      | 3.80 ± 0.5           |        | 3%                   | (Acciari et al. 2010a) |
| H 1426+428            | BBL      | 0.129    | 18.55                                 | 7         | average      | ...                 |        | 3%                   | (Horns et al. 2004) |
| PG 1553+113           | HBL      | 0.4*     | 16.49                                 | 44         | high         | 3.4 ± 0.1           |        | 8%                   | (Aleksić et al. 2010) |
| Markarian 501*        | HBL      | 0.034    | 16.84                                 | 46         | low          | 2.79 ± 0.12         | 20%    |                      | (Anderhub et al. 2009b) |
| iES 1959+650*         | HBL      | 0.048    | 18.03                                 | 16         | low          | 2.58 ± 0.18         |        | 10%                  | (Tagliaferri et al. 2008) |
| PKS 2005-489*         | BBL      | 0.071    | ...                                   | 9         | average      | 3.20 ± 0.16         |        | 3%                   | (Acero et al. 2010) |
| PKS 2155-304*         | HBL      | 0.117    | 15.7                                  | 63         | low          | 3.34 ± 0.05         | 14%    |                      | (Aharonian et al. 2009) |
| BL Lacertae          | LBL      | 0.069    | 14.28                                 | 35         | high         | 3.64 ± 0.54         |        | 3%                   | (Albert et al. 2007c) |
| iES 2344+514*         | HBL      | 0.044    | 16.4                                  | 10         | average      | 2.95 ± 0.12         |        | 11%                  | (Albert et al. 2007a) |
| H 2356-309            | HBL      | 0.165    | 17.24                                 | 8         | average      | 3.06 ± 0.15         |        | 2%                   | (Abramowski et al. 2010a) |

Notes. Columns (1), (2), and (3) show the spectral energy distribution (SED) type, redshift, and synchrotron peak frequency ($\nu_{\text{peak}}$), respectively. Fermi variability indices (4) were taken from the 1FGL catalog (Abdo et al. 2010b). Fermi states (5) are identified in this work using 27-month Fermi light curves as described in Section 5. In cases where Fermi data are contemporaneous with TeV observations, the corresponding TeV instruments are listed in Column (5). TeV spectral indices (6) were taken from the references listed (9). TeV integral fluxes (7) are above the listed energy threshold (8) and in units of Crab Nebula flux. For the Crab Nebula unit conversions, spectral measurements above 350 GeV from Çelik (2008) are used.

* Miller et al. (1978); Lanzetta et al. (1993).
* Prandini et al. (2011).
* Mazin & Goebel (2007).
* Blazars that are reported as variable in the TeV band, according to TeVCat (http://tevcat.uchicago.edu/).

For VHE data that were taken during the Fermi era, time periods of a few months that cover the corresponding VHE observations were selected for the Fermi spectral analysis. For blazars that have VHE spectra measured before the Fermi era, the first 27 months of Fermi data were analyzed (from 2008 August 4 to 2010 November 4). In all the analysis steps, an energy selection from 300 MeV to 100 GeV was applied to the data.

The Fermi data were analyzed in the following way. First, a 27-month light curve analysis was performed for each source using an aperture photometry technique. Diffuse class events from a region of 1° radius from the target location were selected and counts were plotted as a function of time, each time bin containing 49 counts, corresponding to a signal to noise ratio of 7. For sources with high statistics, low- and high- flux states were identified and separated using the average count rate as a threshold. Figure 7 shows the resulting light curves for all sources, with fluxes normalized to arbitrary units. It should be noted that in this analysis, no background subtraction was performed and therefore the resulting light curves merely give an estimate of high- and low-state time slices.

Next, a spectral analysis was done for each data set. Diffuse class events from a region of interest of 8° radius were selected and analyzed with Fermi Science Tools v9r18p6, using instrument response functions P3_V6_DIFFUSE. Sources from the first Fermi-LAT (1FGL) catalog (Abdo et al. 2010b), bright spots with test statistics > 25 and standard galactic and isotropic

7 http://fermi.gsfc.nasa.gov/ssc/data/analysis/scitools/overview.html
Figure 7. Fermi-LAT aperture photometry light curves with no background subtraction, normalized to arbitrary units. The solid lines, representing the average counts per area per time, separate “low” and “high” flux states, that are later on used to produce “low” and “high” state spectra. In the case of Markarian 421 and 3C 279 light curves, the dashed lines represent $1 \sigma$ deviation from the average, dividing the data set into three separate flux states (“low,” “medium,” and “high”). The shaded areas show the contemporaneous time windows with the corresponding TeV instruments.

(A color version of this figure is available in the online journal.)
Figure 7. (Continued)
Figure 7. (Continued)
diffuse emission background components\(^8\) within the region of interest were included in the source model files. Unbinned maximum-likelihood analysis as described in Cash (1979); Mattox et al. (1996) was applied to each data set, assuming a power-law (PL) spectrum as given in Equation (2).

\[
dN/dE = N_0(E/E_0)^{-\Gamma}.
\]

Additionally, to look for possible spectral features in the data, spectral points were calculated and fitted with different power-law functions, and the results were compared. See Section 6.4 for more details.

Finally, combined GeV–TeV SED data sets were constructed using archival TeV spectra and the corresponding flux state information from references shown in Table 1. With each TeV spectrum, the most suitable Fermi data subset (average, low- or high- state) was used for further study.

6. RESULTS AND DISCUSSION

Out of 26 blazars, 12 did not have enough statistics for a temporal separation of the Fermi data set into different flux states. Therefore for this subsample, an average spectrum was calculated using the entire data set. Data from another subsample with 12 blazars were split into high and low-flux states as described in Section 5. Data from the two brightest blazars (Markarian 421 and 3C 279) were split into three subsets, with low, medium, and high-flux states. See Table 2 for a summary of our Fermi data analysis results.

Our analysis results are consistent with the 2FGL catalog (Nolan et al. 2012). We used the combined GeV–TeV SEDs (see Figure 8) to estimate the IC peak frequency band of each blazar (see Section 6.1). Our sample contains a handful of candidate “TeV-peaked” blazars that we discuss in Section 6.2. In addition, considering the fact that Fermi spectral indices do not vary significantly between low- and high-states, we studied the change in spectral index from GeV to TeV as a function of the redshift, thus confirming the EBL effect on TeV spectra with a model-independent approach (see Section 6.3). On the other hand, interesting spectral features in the GeV band are observed. To probe these features, the data were fitted with three different functions and the corresponding fit improvements were calculated (see Section 6.4). Finally, in Section 6.5, we extended this study to contemporaneous combined SEDs.

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\(^8\) http://fermi.gsfc.nasa.gov/ssc/data/access/lat/BackgroundModels.html

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6.1. IC Peak Frequency

The peak frequency of the IC component is a salient parameter for describing blazar non-thermal continua and studying population trends. Systematic studies for measuring the IC peak frequency mostly suffer from the lack of statistics and simultaneous data. A similar work was carried out in Zhang et al. (2012), where archival multiwavelength data were used to model TeV blazar SEDs and determine the IC peak frequency \(\nu_{\text{IC}}\). A positive correlation between \(\nu_{\text{syn}}\) and \(\nu_{\text{IC}}\) was reported.

In this work, we focus on finding the IC “peak frequency band” rather than the “peak frequency,” using a model-independent approach. For each blazar SED shown in Figure 8, we identify the energy decade in which the largest amount of power is emitted. Note that the spectral points used in the VHE spectra are EBL-corrected. Figure 9 shows the distribution of the IC peak bands for different blazar types. We observe that the FSRQs, LBLs and IBLs have the maximum of their emission mostly below 1 GeV. On the other hand, HBLs tend to peak in the TeV range. This positive correlation between the synchrotron \(\nu_{\text{syn}}\) and the IC peak frequencies \(\nu_{\text{IC}}\) is in accordance with simple SSC models that predict a positive correlation between \(\nu_{\text{syn}}\) and \(\nu_{\text{IC}}\) (Abdo et al. 2010a). The dashed lines represent the same distributions with the bright AGN sample from the first three months of Fermi data (Abdo et al. 2010a). Our results tend to span the high frequency sides of all distributions and one clearly sees a shift to higher frequencies in the case of HBLs. This is expected since our sample consists of TeV-selected objects that mostly correspond to relatively weak sources in the GeV data and are therefore less likely to appear in a bright AGN sample. It should also be noted that we use a model independent method using only Fermi and VHE data, whereas (Abdo et al. 2010a) uses multwavelength data and some modeling in cases where the soft X-ray band is dominated by the synchrotron component, a typical feature for our blazar sample.

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6.2. Hard TeV BL Lac Objects

The combined GeV–TeV spectra of some blazars in our sample (1ES 0229+200, 1ES 0347-121, 1ES 1101-232, 1ES 1218+304, H 1426+428) suggest a \(\nu_{\text{IC}}\) beyond ~1 TeV. These blazars are mostly weak or non-detected in the Fermi range, with a hard spectral index in both GeV and TeV bands. It follows that they may belong to the so-called ultra-high-frequency-peaked BL Lac subclass (UHBLs; see, e.g., Costamante 2011) that would constitute the extreme end of the population, and is expected to dominate the TeV luminosity of the universe. Several mechanisms have been set forth to explain the formation
of these hard γ-ray spectra (Lefa et al. 2011). Extensive spectral analysis of these objects would be valuable for EBL and intergalactic magnetic field measurements. It should be noted that at energies of a few TeV and beyond, our spectra become EBL-model-dependent. For this reason, we have compared our adopted EBL model with two other models from recent studies (Finke et al. 2010; Franceschini et al. 2008). We have found that for the data samples mentioned above, if we used any of the other two EBL models, the dispersion in the highest energy flux points would be less than 20%, and consequently the observed spectral upturns would not be affected significantly.

With additional data, a deeper variability study carried on these blazars would relate to arguments that support the cosmic ray production as the origin of TeV blazar emission, since in that scenario no short timescale variability would be expected to be observed (Murase et al. 2012). Among the UHBL candidates, the ones that are present in the 1FGL catalog (1ES 1101-232, 1ES 1218+304, H 1426+428) have relatively small Fermi variability indices (see Table 1). In addition to that, our calculations of $F_{\text{var}}$ for all five blazars using 27 months of Fermi data do not indicate a significant hint of variability either (see Table 2).

### 6.3. Spectral Variability

VHE emission from blazars is highly variable. This variability, manifested in irregular flares, is one of the most typical and promising blazar behaviors for studying the nature of underlying emission mechanisms. The observed flux change during a...
Figure 8. GeV–TeV spectra for the sample of blazars in this study. The filled circles represent the Fermi spectra and the filled squares the TeV spectra. Considering the TeV flux state information given in the TeV papers, the best matching GeV and TeV spectral points are used for the combined analysis (shown in black). When available, spectral points belonging to other flux states (in both bands) are plotted in gray. 3C66A (VERITAS and MAGIC, respectively), RGB J0710+591, 1ES 1218+304, 1222+21 (4C+21.35) (VERITAS), PKS 1424+240, PKS 2155-304 spectra are quasi-simultaneous.
VHE flare can be as rapid as minute scales (Mkn 501: Albert et al. 2007e) and as large as 40 times the baseline emission (Arlen et al. 2013). Blazars that have been reported to have a variable flux are marked with an asterisk in Table 1. On the other hand, Fermi data do not exhibit flux variability as extreme as in the VHE band. In fact, having a smaller effective area than the ground-based VHE telescopes and operating mostly in survey mode rather than pointing, Fermi-LAT does not have the sensitivity to probe sub-hour timescale variability in blazars. Still, a possible correlation between GeV and TeV emission

Figure 8. (Continued)
remains viable (Abdo et al. 2011; Aleksic et al. 2011b) and an enhanced activity in the high-energy tail of the Fermi band could therefore indicate a TeV flare. In this frame, monitoring GeV flares to trigger TeV observations is important (Errando et al. 2011a) and potentially could help in probing fast variability. To examine variability within the Fermi data, we compared high- and low-state Fermi spectra from 14 blazars (see Table 3). Half of these blazars have their integral flux in 1–100 GeV
Notes. Column (1) shows the spectral energy distribution (SED) type. $F_{\gamma-100}$ is the integral flux for the 1–100 GeV band. Column (2) shows the percent increase in $F_{\gamma-100}$ from low to high Fermi state. FSRQs and LBLs seem to show the most significant flux variability in this energy range. Columns (3) and (4) list the GeV photon indices in low and high Fermi states, respectively. No significant changes in photon index are seen, except for the two FSRQs for which the index shows a slight hardening from low to high flux states.

Table 3

| Name             | SED Type | Increase in $F_{\gamma-100}$ (%) | $\Gamma_{\text{low}}$ (3) | $\Gamma_{\text{high}}$ (4) |
|------------------|----------|----------------------------------|--------------------------|--------------------------|
| 3C 66A           | BLM      | 100                              | 2.10 ± 0.02              | 1.93 ± 0.02              |
| S5 0716+714      | LBL      | 100                              | 2.10 ± 0.02              | 1.30 ± 0.04              |
| 1ES 1011+496     | HBL      | 30                               | 1.90 ± 0.04              | 1.97 ± 0.04              |
| Mrk 421          | HBL      | 100                              | 1.86 ± 0.03              | 1.79 ± 0.02              |
| 1ES 1218+304     | HBL      | 100                              | 1.68 ± 0.08              | 1.73 ± 0.10              |
| W Comae          | HBL      | 100                              | 2.10 ± 0.06              | 2.07 ± 0.06              |
| PKS 1222+21      | FSRQ     | 1290                             | 2.50 ± 0.04              | 2.32 ± 0.02              |
| 3C 279           | FSRQ     | 295                              | 2.54 ± 0.03              | 2.37 ± 0.02              |
| PKS 1424+240     | BLM      | 45                               | 1.82 ± 0.04              | 1.87 ± 0.03              |
| PG 1553+113      | HBL      | 30                               | 1.70 ± 0.03              | 1.74 ± 0.03              |
| Mrk 501          | HBL      | 100                              | 1.84 ± 0.05              | 1.83 ± 0.04              |
| 1ES 1959+650     | HBL      | 100                              | 2.04 ± 0.06              | 2.07 ± 0.05              |
| PKS 2155-304     | BLM      | 80                               | 1.95 ± 0.03              | 1.92 ± 0.02              |
| BL Lacertae      | LBL      | 140                              | 2.46 ± 0.05              | 2.34 ± 0.04              |

(Continued)
with the EBL photons (Franceschini et al. 2008) and this effect becomes more enhanced at larger redshifts, making the universe opaque to TeV $\gamma$ rays at distances larger than $z \sim 0.5$. HE spectra are not affected by the EBL, whereas VHE spectra become softer with increasing redshift. A similar observation was reported by Abdo et al. (2009), in a study carried out on a sample of TeV-selected AGNs detected with Fermi.

Figure 11 shows the relation between the spectral index $\Gamma_{\text{GeV}}$ and the flux normalization $F_{1-100}$ obtained from power-law fits. FSRQs and two subgroups of BL Lacs are clearly separated in the parameter space. This is in accordance with the aforementioned positive correlation trend between $v_{\text{syn}}$ and $v_{\text{IC}}$, since 1 GeV typically corresponds to the rising edge of the IC component in an HBL SED, sampling a relatively low flux with hard spectral index. On the other hand for an FSRQ, 1 GeV will correspond to the peak or the falling edge of the IC component. The fact that FSRQs have relatively more luminous IC emission explains the softening trend with a larger normalization factor. However, the pattern that we observe between different flux states of a given blazar is the opposite. In most cases, a slight spectral hardening accompanies high flux states, indicating a change in the spectral shape and enhanced flux increase at high-energy tail of the spectrum.

6.4. Spectral Features

In most of the blazars in our sample, we observe interesting spectral features in the Fermi band, that appear as dips in the 1–100 GeV energy range. In an attempt to find a quantitative description for these features, we fit the Fermi spectral points with a simple power law (PL; Equation (2)) and a broken power law (BPL; Equation (3)), and then compare the results:

$$dN/dE = N_0 \times \begin{cases} (E/E_b)^{-\Gamma_1}, & E < E_b \\ (E/E_b)^{-\Gamma_2}, & \text{otherwise.} \end{cases}$$

(3)

In the PL fit, the normalization $N_0$ and the spectral index $\Gamma$ are free parameters, and the energy $E_b$ is fixed at 1 GeV. In the BPL fit, the break energy $E_b$ and the indices $\Gamma_1$ and $\Gamma_2$, along with the normalization $N_0$ are free. In Table 4, we list the best-fit parameters from both functions and the likelihood ratio test results of BPL over PL. In 9 out of 33 cases, BPL yields a better fit over PL with more than $2\sigma$ significance.

There are several possible mechanisms that may cause the observed features in the SEDs. One possibility is a break in the electron spectrum caused by the synchrotron cooling effects, generally yielding a change in spectral index by 0.5 (Böttcher & Chiang 2002), which is in agreement with our results (see Table 4). Another mechanism that could explain the observed
breaks is the absorption by an external photon field (Poutanen & Stern 2010). For those nine data sets where the BPL gives a better fit than the PL, the break energy ranges from ~2 GeV to ~8 GeV. In addition, seven of these data sets belong to non-HBL blazars, that are usually characterized by broad emission lines, thought to be originating from a region of molecular gas (broad line region, BLR) that is highly ionized by the optically thin accretion disk. This seems in accordance with the idea of relating the Fermi spectral features to absorption of GeV photons on radiation from H I (13.6 eV) and He II (54.4 eV) recombination continua in the BLR, that are expected to cause jumps in γ-ray opacity around ~19.2 and ~4.8 GeV, respectively (Poutanen & Stern 2010). We tested a general absorbed power-law (APL) function of the following form on the Fermi data:

\[ \frac{dN}{dE} = N_0 (E/E_0)^{-\Gamma} e^{-\tau_{\gamma\gamma}(E,z,E_{\text{abs}})}, \]

where the free parameters are the normalization \( N_0 \) at \( E_0 = 1 \) GeV, photon index \( \Gamma \), and absorption line energy \( E_{\text{abs}} \). \( \tau_{\gamma\gamma} \) is the optical depth for the \( \gamma-\gamma \) pair annihilation of photons with energies \( E \) and \( E_{\text{abs}} \) at a redshift of \( z \). Within the Fermi energy band, BPL and APL functions fit the data equally well. Upturns at high-energy tails of Fermi spectra are observed (see, e.g., W Comae in Figure 8), but they are not statistically significant enough to favor an absorption scenario over a BPL fit. Therefore, it appears that one cannot statistically distinguish between the BPL and APL fits, but possible absorption scenarios are worth investigating further. To address this issue, we make use of contemporaneous GeV–TeV spectra to test and compare BPL and four different APL scenarios (see Section 6.5). This permits us to test the APL over a larger energy range and investigate the apparent Fermi spectral absorption-like features with higher statistics.

Another caveat related to these spectral features is that the upturn seen at the highest Fermi energy bin might be coming from a group of photons clustered in time. In that case the dip would be an artifact of a flaring event, thus not representative of the time-averaged spectrum. To make sure this is not the case, we checked the arrival times of the highest-energy photons and did not find any obvious clustering (see Figure 12). Note that the arrival time distributions should be considered within a given flux state. For instance, in the left panel of the figure, the red triangles represent the high energy photons from the high flux state and are evenly distributed in a time window that belongs to the high state. Therefore, one concludes that no clustering is found.

6.5. Quasi-simultaneous GeV–TeV Spectra

Seven of the TeV spectra in our sample are contemporaneous with Fermi observations and therefore merit a deeper analysis. We extended the work described in Section 6.4 to this subsample. This time, in addition to PL and BPL fits, we tested four different scenarios of absorption due to photons emitted from the BLR: H I line (13.6 eV), He II line (54.4 eV), H I and He II combined, and full BLR spectrum taken from Poutanen & Stern (2010).

For single- and double-line absorption scenarios, H I and He II recombination continua are the most plausible cases given that they are the most dominant ones in the BLR spectrum, and that the breaks we see in the Fermi spectra are located around a few GeV. As for the full BLR spectrum, taken from Poutanen & Stern (2010), it was modeled assuming a photoionized gas with the ionization parameter and the cloud density changing with the distance to the central ionizing source. See Poutanen & Stern (2010) for a detailed discussion on the γ-ray absorption within the BLR in the Fermi spectra. No general trend can be seen in the contemporaneous data sample. BPL and full BLR absorption scenarios seem to fit well the combined spectra of the blazars PKS 1424+240 and PKS 2155-304. A BPL (full BLR absorption) function is preferred over the PL for the blazars PKS 1424+240 and PKS 2155-304 with a significance of ~3σ (~4.8σ) and ~12σ (~8.5σ), respectively (see Figure 13). The χ²/dof values of PL, BPL, and APL fits are 32/9, 3.5/7, and 5.3/7 for PKS 1424+240 and 148/8, 5.8/6, and 71/6 for PKS 2155-304, respectively. BPL fits yield ΔΓ = 1.4 (PKS 1424+240) and ΔΓ = 0.7 (PKS 2155-304), both larger than what electron cooling would predict, which might indicate that an additional mechanism is at work. Both BPL and full BLR absorption scenarios provide a slight improvement in the MAGIC and VERITAS spectra of 3C 66A, albeit not significant. Similarly, for PKS 1222+21, BPL, H I single line and H I + He II double line absorptions slightly improve the fit over PL. In the case of RGB J0710+091 and IES 1218+304, we don’t observe any preference over the power-law fit. In case a γ–γ absorption from BLR is at work, the cascades initiated in this process might produce observable GeV γ-ray emission (Roustazadeh & Böttcher 2011), and their synchrotron emission could contribute

![Table 4](https://example.com/table4.png)

### Table 4

**Fit Results for Power Law (PL) and Broken Power Law (BPL), Where BPL Yields a Better Fit Over PL with More Than 2σ Significance (9 Out of 33 Cases)**

| Name              | N(x10^{-11}) | Power Law | Broken Power Law |
|-------------------|--------------|-----------|-----------------|
|                   | (1)          | \( \Gamma \) | \( \Gamma_1 \) | \( \Gamma_2 \) | \( E_{\text{break}} \) (GeV) | \( \sigma_{\text{BPL}} \) (7) |
| 3C 279 (low)      | 2.78 ± 0.07  | 2.01 ± 0.02 | 9.30 ± 4.94 | 1.81 ± 0.06 | 2.22 ± 0.07 | 1.92 ± 0.51 | 4.12 |
| 3S 0716+714 (low) | 1.85 ± 0.06  | 2.23 ± 0.03 | 5.18 ± 2.07 | 2.14 ± 0.04 | 2.74 ± 0.23 | 5.43 ± 1.01 | 3.07 |
| PKS 2155-304     | 2.96 ± 0.13  | 2.10 ± 0.03 | 2.29 ± 1.78 | 2.00 ± 0.06 | 2.57 ± 0.30 | 4.10 ± 1.47 | 2.29 |
| 3ES 1222+46 (low) | 0.88 ± 0.04  | 1.94 ± 0.04 | 1.70 ± 1.34 | 1.70 ± 0.10 | 2.28 ± 0.20 | 2.76 ± 1.15 | 2.76 |
| PKS 2155-304 (low) | 1.72 ± 0.06  | 2.53 ± 0.04 | 0.99 ± 0.70 | 2.41 ± 0.06 | 3.11 ± 0.28 | 3.36 ± 0.87 | 2.84 |
| 3C 279 (low)      | 2.55 ± 0.08  | 2.59 ± 0.03 | 0.15 ± 0.28 | 2.53 ± 0.04 | 3.72 ± 1.54 | 7.70 ± 5.57 | 2.51 |
| 3C 279 (high)     | 1.16 ± 0.02  | 2.39 ± 0.02 | 8.14 ± 10.72 | 2.31 ± 0.03 | 2.74 ± 0.26 | 3.21 ± 1.79 | 3.11 |
| PKS 2155-304 (low) | 4.19 ± 0.10  | 1.92 ± 0.02 | 1.75 ± 1.98 | 1.86 ± 0.03 | 2.13 ± 0.14 | 5.54 ± 3.23 | 2.08 |
| BL Lacertae (high)| 3.26 ± 0.11  | 2.39 ± 0.04 | 1.08 ± 0.09 | 2.30 ± 0.04 | 3.00 ± 0.30 | 4.94 ± 0.03 | 2.50 |

**Notes.** Columns (1) and (2) show the PL parameters, flux normalization at 1 GeV and the photon index respectively. \( F_{\text{GeV}} \) (3) is the flux normalization at 1 GeV for BPL. \( N \) and \( F_{\text{GeV}} \) are in erg cm⁻² s⁻¹. Columns (4) and (5) show the photon indices for BPL, as given in Equation (3). The break energy for BPL is listed in Column (6). \( \sigma_{\text{BPL}} \) is the likelihood ratio test results of BPL over PL. For these 9 cases, the break energy \( E_{\text{break}} \) ranges from ~2 GeV to ~8 GeV. In addition, 7 of these data sets belong to non-HBL blazars.
Figure 12. Arrival times of the photons in the highest energy bin (top panels) and aperture light curves with arbitrary flux units (bottom panels) for the blazars 3C 279 (left) and PKS 2155-304 (right). The blue upside-down triangles represent the low-state photons and red triangles the high-state ones. In both cases, the highest energy photons do not show any obvious clustering within their respective data sets. As described in Section 5, low- and high-states are distinguished based on the flux averages (solid lines) in light curves.

(A color version of this figure is available in the online journal.)

Figure 13. Contemporaneous GeV–TeV spectra with power-law (dashed lines), broken power-law (solid lines), and power-law with full-BLR-absorption fits. BPL and full BLR absorption scenarios seem to fit well the combined spectra of the blazars PKS 1424+240 and PKS 2155-304. A BPL (full BLR absorption) function is preferred over the PL for the blazars PKS 1424+240 and PKS 2155-304 with a significance of $\sim 5\sigma$ ($\sim 4.8\sigma$) and $\sim 12\sigma$ ($\sim 8.5\sigma$), respectively.

(A color version of this figure is available in the online journal.)
to the big blue bump seen in several blazars (Roustazadeh & Böttcher 2012).

7. SUMMARY

We study blazar spectral properties with a focus on the GeV–TeV energy range for a sample of VHE blazars. In order to obtain a set of joint GeV–TeV blazar spectra, we analyze the first 27-month Fermi data for VHE blazars and combine our results with archival VHE data. In cases where the Fermi data set does not overlap with the TeV observations but has enough statistics, we split the data into high and low flux states and join the best-matching subset with the corresponding TeV spectrum. The peak frequency band of the inverse Compton component increases following the order FSRQ -> LBL&IBL -> HBL. Thus, our results confirm the positive correlation between \( \nu_{\text{syn}} \) and \( \nu_{C} \). We note that Fermi spectra from different flux states for a given TeV blazar do not undergo a significant change in photon index. The variability amplitudes within our Fermi data set do not show an immediate correlation with the reported TeV variabilities for individual blazars. We find that in many cases a power law is insufficient to describe the GeV–TeV spectra and a broken power law improves the fits, especially for non-HBL blazars, where the BLR emission may have an effect on the observed spectral shape. In some blazars we observe absorption-like spectral features. We present seven quasi-simultaneous joint spectra, on which we test possible absorption scenarios from the BLR. Even though the absorption seems to describe well some of the observed spectra, no general pattern can be identified.

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