Brightening X-Ray Emission from GW170817/GRB 170817A: Further Evidence for an Outflow

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Received 2017 December 7; revised 2017 December 23; accepted 2018 January 2; published 2018 January 18

Abstract

The origin of the X-ray emission from neutron star coalescence GW170817/GRB 170817A is a key diagnostic of the unsettled post-merger narrative, and different scenarios predict distinct evolution in its X-ray light curve. Due to its sky proximity to the Sun, sensitive X-ray monitoring of GW170817/GRB 170817A has not been possible since a previous detection at 16 days post-burst. We present new, deep Chandra observations of GW170817/GRB 170817A at 109 days post-burst, immediately after Sun constraints were lifted. The X-ray emission has brightened from a 0.3–8.0 keV flux of $3.6 \times 10^{-15}$ erg s$^{-1}$ cm$^{-2}$ at 16 days to $15.8 \times 10^{-15}$ erg s$^{-1}$ cm$^{-2}$ at 109 days, at a rate similar to the radio observations. This confirms that the X-ray and radio emission have a common origin. We show that the X-ray light curve is consistent with models of outflow afterglows, in which the outflow can be a cocoon shocked by the jet, dynamical ejecta from the merger, or an off-axis structured jet. Further deep X-ray monitoring can place powerful constraints on the physical parameters of these models, by both timing the passing of a synchrotron cooling break through the X-ray band and detecting the associated steepening of the X-ray photon index. Finally, the X-ray brightening strengthens the argument that simple off-axis top-hat jet models are not consistent with the latest observations of GW170817/GRB 170817A.

Key words: gamma-ray burst: individual (GRB 170817A) – gravitational waves – stars: neutron – X-rays: binaries

1. Introduction

The gravitational-wave (GW) and multi-wavelength electromagnetic (EM) discoveries of the binary neutron star (NS) merger GW170817 marked the dawn of multi-messenger GW astronomy (e.g., Abbott et al. 2017a, 2017b; Coulter et al. 2017; Evans et al. 2017; Goldstein et al. 2017; Hallinan et al. 2017; Soares-Santos et al. 2017; Troja et al. 2017; Valenti et al. 2017). Detection of the short gamma-ray burst (sGRB) GRB 170817A associated with the gravitational-wave event GW170817 confirmed that binary NS mergers are the progenitors of at least some sGRBs (Abbott et al. 2017c; Goldstein et al. 2017; Savchenko et al. 2017). Furthermore, observations of the optical and infrared transient confirmed that binary NS mergers produce kilonova, powered by r-process nucleosynthesis (Arcavi et al. 2017; Cowperthwaite et al. 2017; Drout et al. 2017; Kasliwal et al. 2017; McCully et al. 2017; Pian et al. 2017; Shappee et al. 2017; Smartt et al. 2017). Finally, robust association of GW170817 with its host galaxy (Hjorth et al. 2017; Im et al. 2017; Levan et al. 2017) enabled measurements of the Hubble constant independent of the cosmic distance ladder (Abbott et al. 2017d; Guidorzi et al. 2017), providing a new probe of cosmology. However, many questions about the nature of GW170817 remain open.

X-ray observations of GW170817 provided important early indication that the associated sGRB GRB 170817A was unlike any other GRB afterglow previously observed (Fong et al. 2017), Chandra observations of the optical counterpart (SSS17a; Coulter et al. 2017) at ~2 days post-burst yielded a non-detection (Margutti et al. 2017). However, Chandra observations at ~9 days post-burst revealed a new X-ray source coincident with the position of the optical transient (Troja et al. 2017). This delayed brightening of the X-ray counterpart has not been observed in standard GRB afterglows, which instead display consistent dimming in their X-ray light curves on timescales of days. The discovery of an X-ray counterpart was confirmed in additional Chandra observations at ~15 and 16 days post-burst, which surprisingly showed that the source had not dimmed significantly (Haggard et al. 2017). Unfortunately, Sun constraints prevented further X-ray monitoring of GRB 170817A after the last detection at 16 days post-burst, until early 2017 December (~109 days post-burst).

The origin of the observed X-ray emission from GRB 170817A is still unclear. Haggard et al. (2017), Margutti et al. (2017), and Troja et al. (2017) all showed that the Chandra X-ray light curve is consistent with the fading of a standard sGRB afterglow in which a simple top-hat jet is off-axis from the line of sight by $20^{\circ}$–$30^{\circ}$. This interpretation is ostensibly supported by the low fluence observed in the prompt $\gamma$-ray emission (e.g., Abbott et al. 2017c; Goldstein et al. 2017; Murguia-Berthier et al. 2017), which was a factor of $\sim 10^{3}$ lower than previously observed in any other sGRBs. Modeling of the radio light curves of GRB 170817A similarly showed consistency with off-axis sGRB afterglows (Alexander et al. 2017; Hallinan et al. 2017). If confirmed, this also makes
In this Letter, we present deep Chandra observations of GRB 170817A at 109.2 days post-burst. Due to Sun constraints, these are the first X-ray observations since the previous Chandra detection at 16 days post-burst. The outline of this Letter is as follows: In Section 2, we describe our Chandra observations and data reduction procedure. In Section 3, we discuss the origin of the X-ray and radio emission, and compare the observed X-ray light curve to predictions from outflow and structured jet afterglow models. We briefly summarize and conclude in Section 4.

2. X-Ray Observations and Analysis

New X-ray observations of GRB 170817A were obtained via a Chandra Director’s Discretionary Time allocation (PI: Wilkes, Program Number 18408601). This program obtained two exposures of GRB 170817A: (1) a 74.09 ks exposure (ObsID 20860) beginning at 2017 December 2.08 UT, approximately 108 days post-burst, and (2) a 24.74 ks exposure (ObsID 20861) beginning at 2017 December 6.45 UT, approximately 111 days post-burst. Both these exposures were acquired using Chandra’s ACIS-S3 chip in VFAINT mode. We use CIAO v.4.9.4 (CALDB v4.7.6; Fruscione et al. 2006) for reduction and analysis of these X-ray data. We reprocess the level 2 event files and use CIAO’s reproc script to apply the latest calibrations. Since the two new exposures are close in time and the X-ray emission of GRB 170817A is not expected to vary significantly over ∼4 day timescales, we co-add the two data sets into one 98.83 ks exposure at 109.2 days post-burst.

Figure 1 shows our latest Chandra 0.5–8 keV image of GRB 170817A at 109.2 days post-burst (right), in comparison to the Chandra image from a similar exposure at 15.6 days post-burst from Haggard et al. (2017) (left). The previously detected X-ray source at the position of GRB 170817A is still detected in this latest observation, along with the three other previously detected X-ray sources in the field: CXOU J130948, CXOU 130946, and the host galaxy NGC 4993. We construct a difference image by subtracting the 15.6 day image from the 109.2 day image, scaled by their exposure times. This difference image (Figure 1, right) shows that GRB 170817A has brightened in the latest observations.

To measure the flux of each of the four sources, we first apply the point-source detection algorithm wavdetect to each 0.5–7 keV image to determine the centroid positions of each source. Following the same procedure from Haggard et al. (2017), we extract spectra from the point sources GRB 170817A, CXOU J130948, and CXOU 130946 using extraction regions with radii 1.9′ (corresponding to ∼90% encircled energy fraction near the Chandra on-axis position). For the host galaxy NGC 4993 we adopt an extraction region with a 2.9′ radius, large enough to include most of the galaxy’s X-ray flux while avoiding contamination from nearby CXOU J130948. These regions are shown in Figure 1. We obtain background photons from a large region on the same chip that does not overlap other sources.

We extract spectra and response files for the four detected X-ray sources using the CIAO specextract tool from the individual observations. The resultant files are then co-added to improve statistics. We fit the combined spectra using XSPEC v12.9.0 (Arnaud 1996), with atomic cross sections from Verner et al. (1996) and abundances from Wilms et al. (2000). For each of the four sources, we assume absorbed power-law spectral models with fixed $N_H = 7.5 \times 10^{20} \text{cm}^{-2}$. The best-fit
We test for additional model components by adding Gaussians at various line centers to the absorbed power-law model. Though adding a line component, e.g., at $E = 1.8$ keV, can lower the reduced $\chi^2$ to 0.49 and remove the excess around 2 keV, the Gaussian does not dominate the spectrum. The spectral index is also reduced, but its larger uncertainty ($\Gamma = 1.45^{+0.45}_{-0.36}$) places the value within the error interval of the power-law-only fit. Hence, inclusion of an emission line is not yet statistically supported by the *Chandra* data. If GRB 170817A continues to brighten we may be able to better quantify this tantalizing line emission. We also note that a simple absorbed thermal model (blackbody, $kT = 0.63 \pm 0.09$) can fit to the spectrum with a similar reduced $\chi^2$. However, none of the current models for the late-time X-ray emission predict a thermal spectrum, and the physical process that links the brightening radio and X-ray emission would be unclear.

The flux and spectrum of the host galaxy NGC 4993 and CXOU J130948 are consistent with our previous deep *Chandra* observations. CXOU J130948 is known to be variable in X-ray, and while our spectral analysis shows its spectrum is consistent with those previously reported, the best-fit parameters are not

### Table 1

| Source ID | Days Post-burst | Count Rate (0.5–10 keV) ($10^{-3}$ cts s$^{-1}$) | Power Law Index $\Gamma$ | Flux (0.3–8 keV) ($10^{-16}$ erg s$^{-1}$ cm$^{-2}$) | Luminosity ($10^{38}$ erg s$^{-1}$) | $\chi^2$/dof |
|-----------|----------------|---------------------------------|-------------------|-------------------|----------------|----------------|
| GRB 170817A | 15.6 | 3.5 ± 0.7 | $2.4_{-0.8}^{+0.8}$ | $0.36_{-0.10}^{+0.07}$ | $10.4_{-1.9}^{+2.0}$ | 3.15/5 |
| | 109.2 | 14.7 ± 1.3 | $1.6_{-0.3}^{+0.2}$ | $1.5_{-0.11}^{+0.14}$ | $42.5_{-3.3}^{+17.2}$ | 11.9/13 |
| CXOU J130948 | 15.6 | 2.5 ± 0.6 | $1.3_{-0.8}^{+0.8}$ | $0.4_{-0.09}^{+0.10}$ | $10.8_{-2.4}^{+4.4}$ | 1.77/3 |
| | 109.2 | 1.2 ± 0.4 | $2.4_{-1.7}^{+1.7}$ | $0.2_{-0.11}^{+0.13}$ | $6.5_{-0.7}^{+4.2}$ | 0.26/2 |
| CXOU J130946 | 15.6 | 3.3 ± 0.7 | $-0.4_{-0.8}^{+0.2}$ | $1.1_{-0.1}^{+0.1}$ | $38.8_{-9.8}^{+4.0}$ | 3.68/5 |
| | 109.2 | 2.5 ± 0.7 | $0.4_{-0.9}^{+0.9}$ | $0.5_{-0.13}^{+0.15}$ | $17.5_{-10.3}^{+11.9}$ | 2.41/3 |
| NGC 4993 | 15.6 | 11.0 ± 1.2 | $1.5_{-0.4}^{+0.4}$ | $1.3_{-0.2}^{+0.3}$ | $34.8_{-13.2}^{+13.2}$ | 12.73/11 |
| | 109.2 | 12.3 ± 1.2 | $1.4_{-0.3}^{+0.4}$ | $1.4_{-0.3}^{+0.3}$ | $38.9_{-3.4}^{+5.9}$ | 13.09/11 |

Notes. All reported errors represent 90% confidence intervals. The neutral hydrogen absorption was frozen to $N_H = 7.5 \times 10^{20}$ cm$^{-2}$ for all spectral fits, based on NGC 4993’s $A_V = 0.338$ (Schlaufy & Finkbeiner 2011).

a A luminosity distance of 42.5 Mpc was assumed for all sources.

b The luminosities at 15.6 days listed here are larger than reported in Haggard et al. (2017) by a factor of 4, due to a previous calculation error.
Spectral data and residuals are shown in black and the best-fit spectral models are in red. The neutral hydrogen absorption column is fixed to $N_H = 7.5 \times 10^{20}$ cm$^{-2}$ (see Section 2). These spectra at 15.6 and 109.2 days each consist of two exposures that were jointly fitted between the exposures and the power-law normalization was left free.

well constrained. This source is visibly dimmer at 109.2 days (Figure 1) and shows a $\sim 48\%$ decrease in count rate between the two epochs (Table 1; Haggard et al. 2017; Margutti et al. 2017).

3. Discussion

The observed X-ray brightening of GRB 170817A at 109.2 days in comparison to previous X-ray observations at 9 and 15.6 days has profound implications for our understanding of the EM transient. In Section 3.1, we first test whether the X-ray and radio emission have the same origin by comparing the X-ray fluxes to the radio observations. We find that the X-rays brightened at a rate similar to the radio emission, confirming that they share a common origin. In Section 3.2, we consider the possibility that both the X-ray and radio emission are due to the afterglow of an outflow, either from a cocoon, dynamical ejecta, or a structured jet. We find that the X-ray light curve is well described by predictions from outflow models based on radio observations.

3.1. Comparison of X-Ray and Radio Emission: Evidence for a Common Origin

The similar brightening of both X-ray and radio emission from GRB 170817A suggests that the X-ray and radio emission may have a common origin. This can be confirmed if the X-ray-to-radio spectral index $\alpha$ (where $S_{\nu} \sim \nu^\alpha$) is measured to be constant between 15.6 and 109.2 days post-burst. In this scenario, the X-ray light curve would be consistent with a scenario in which afterglow emission from an outflow is the origin of both the radio and X-ray emission, while off-axis GRB afterglows from simple top-hat jet models are ruled out. If instead $\alpha$ is measured to have steepened, this can either indicate that the X-ray and radio emission have different origins, or that they have the same synchrotron origin but a cooling break has passed through the X-ray band.

We compare the X-ray-to-radio spectral index $\alpha$ at 15.6 and 109.2 days and find that it remains constant, confirming a common origin for both the X-ray and radio emission. For the X-ray, we use the measured X-ray flux densities at 15.6 days from Haggard et al. (2017) and at 109.2 days reported here. For the radio, we use the VLA 3 GHz radio light curve from Hallinan et al. (2017) and Mooley et al. (2017), which covers a timespan from 16.4 to 93.1 days post-burst. Mooley et al. (2017) showed that this radio light curve is well fit by a power law with slope $\Gamma = 0.8$, and thus we also fit a power law to extrapolate the radio flux density at 15.6 and 109.2 days. The uncertainties on these two extrapolated radio flux densities are estimated through Monte Carlo re-sampling of the radio flux density measurements in the light curve based on their uncertainties, then re-performing the fit and extrapolation. Figure 3 compares the X-ray and radio flux densities at the two epochs. The fitted X-ray-to-radio spectral index is $\sim -0.53 \pm 0.02$ at 15.6 days and $-0.53 \pm 0.01$ at 109.2 days. Thus, we detect no change in $\alpha$, implying that the X-rays brightened at the same rate as the radio. The similar rate of brightening in the X-ray and radio light curves thus confirms a common origin. Moreover, the consistency between the radio spectral index ($\alpha = 0.60$; Alexander et al. 2017; Mooley et al. 2017) and

![Figure 2](https://example.com/figure2.png) **Figure 2.** Chandra co-added X-ray spectra from GRB 170817A in a 93.4 ks exposure at 15.6 days (left) and a 98.8 ks exposure at 109.2 days post-burst (right). The observed X-ray brightening of GRB 170817A at 109.2 days has profound implications for our understanding of the EM transient. In Section 3.1, we first test whether the X-ray and radio emission have the same origin by comparing the X-ray fluxes to the radio observations. We find that the X-rays brightened at a rate similar to the radio emission, confirming that they share a common origin.

![Figure 3](https://example.com/figure3.png) **Figure 3.** Comparison of the radio to X-ray flux densities at 15.6 days (blue points) and 109.2 days post-burst (red points). The X-ray-to-radio spectral index $\alpha$ remains constant at $-0.53$ over this timespan, indicating that the X-ray emission brightened at the same rate as the radio.
Figure 4. Chandra X-ray light curve of GRB 170817A (black points), including our new observations at 109.2 days post-burst. The predicted X-ray light curves from the cocoon model of Gottlieb et al. (2017; solid orange line), the cocoon (solid red line) and ejecta (solid blue line) outflow models of Mooley et al. (2017), and the structured jet model of Lazzati et al. (2017) are also shown. The gray shaded region is the timespan over which Chandra observations were not possible due to Sun constraints. All uncertainties shown are 90% confidence intervals. The brightening X-ray light curve is well described by cocoon, dynamical ejecta, or structured jet models, although several of these models underpredict the detected early X-ray flux at 9 days post-burst.

X-ray photon index ($\Gamma = 1.62$, where $\alpha = \Gamma - 1$; see Table 1) further supports this conclusion.

3.2. Comparison with Models

Mooley et al. (2017) showed that the brightening radio light curve of GRB 170817A rules out simple top-hat off-axis jet models and is instead well described by models of afterglow emission from an outflow. This outflow can be a mildly relativistic cocoon shocked by the jet head or the high-velocity tail of dynamical ejecta from the neutron star coalescence. In both cases, the slow and monotonic rise of the radio light curve implies that the blast wave must have a continued injection of kinetic energy. Mooley et al. (2017) show that both cocoon and dynamical ejecta models in which the majority of the kinetic energy is in the lower-velocity material provide excellent fits to the radio light curve. These models also predict that the X-ray light curve will rise at the same rate as the radio, which we can directly test with our X-ray light curve.

We produce predicted X-ray light curves based on the radio light curves from both the cocoon and ejecta outflow models presented in Mooley et al. (2017), as well as the cocoon outflow model from Gottlieb et al. (2017). The X-ray emission is estimated by scaling the model radio light curves to the X-rays based on the X-ray-to-radio flux ratio at 15.6 days post-burst. This assumes that the synchrotron cooling break has not yet passed through the X-rays, and thus the X-ray-to-radio spectral index remains constant. This assumption is justified because the fiducial models used by Mooley et al. (2017) to fit the radio light curve predict that the synchrotron cooling frequency is still well above the Chandra band at 109.2 days post-burst. The predicted X-ray light curves from these outflow models are shown in Figure 4.

The X-ray light curves predicted by the three cocoon and dynamical ejecta outflow models in Figure 4 are a good match to our observed X-ray flux at 109.2 days post-burst. Although the cocoon and dynamical ejecta outflow models of Mooley et al. (2017) underpredict the early-time X-ray flux at 9 days, they are fitted to the radio light curve (which do not cover such early times) rather than the X-ray light curve. A more complete fit of these models that also incorporates the X-ray data will likely improve the fit. The general agreement of the X-ray light curve with these models is consistent with a scenario in which the afterglow of a cocoon or dynamical ejecta produce both the X-ray and radio emission and implies that the synchrotron cooling break has not yet passed through the X-ray band.

If the synchrotron cooling frequency is still above the X-ray band, continued deep Chandra observations will provide the first indication of the cooling break, via a change in the X-ray photon index $\Gamma$. Across the break frequency, the X-ray photon index should be observed to steepen by $\Delta \Gamma = 0.5$ (where $\alpha = \Gamma - 1$), from our measured $\Gamma = 1.62 \pm 0.27$ at 109.2 days. Since the synchrotron cooling frequency is strongly dependent on the velocity of the outflow, the detection of a cooling break in X-ray monitoring can provide powerful constraints on critical parameters in these models. We thus strongly encourage additional deep Chandra X-ray observations of GRB 170817A toward this end.

Although Mooley et al. (2017) showed that the radio light curve rules out afterglow emission from off-axis top-hat jet models, more sophisticated off-axis structured jets can still be supported by the current data. For example, Lazzati et al. (2017b) presented a simulation of an off-axis structured jet, and showed that the afterglow emission can produce the observed X-ray and radio properties. When observed off-axis, the afterglow emission from structured jets will come from increasingly more relativistic material closer to the jet axis as the effects of beaming become less severe over time. This can produce a slower brightening in its afterglow light curve in comparison to off-axis top-hat jets, akin to the observed slowly brightening radio and X-ray light curves of GRB 170817A. Figure 4 shows that the predicted X-ray light curve from this off-axis structured jet simulation of Lazzati et al. (2017b) is also well matched to our observed X-ray flux at late times. Since it is currently difficult to robustly differentiate whether the origin of the radio and X-ray emission is from the afterglow...
of a cocoon, dynamical ejecta, or a structured jet, continued X-ray and radio monitoring of GRB 170817A should be avidly pursued.

4. Conclusion

We present late-time Chandra observations of neutron star coalescence GW170817/GRB 170817A at 109.2 days post-burst, the first sensitive X-ray observations possible since a previous detection at 15.6 days. These data show that the X-ray counterpart has brightened at the same rate as its radio light curve. We show that the X-ray light curve is a good match to predictions from outflow models, in which the outflow is a cocoon, dynamical ejecta, or a structured jet. Our observations thus support a scenario in which both the X-ray and radio emission are the afterglow of an outflow, although the exact origin of the outflow is still uncertain. Finally, the X-ray brightening strengthens the argument that simple top-hat jet models are not consistent with the latest observations.

Continued Chandra monitoring of GRB 170817A will be critical for validating outflow models. Since the synchrotron cooling frequency is still above the X-ray band, additional deep observations by Chandra will provide the first indication of the cooling break by detecting a steepening in the X-ray photon index. A detection of the cooling break will provide powerful constraints on the physical parameters of the outflow producing the X-ray and radio emission. Our observations of GW170817/GRB 170817A presented here highlight the unique role of Chandra in opening the multi-messenger gravitational-wave era.

The authors thank Belinda Wilkes and the Chandra scheduling, data processing, and archive teams for making these observations possible. This work was supported by Chandra Award Number GO7-18033X, issued by the Chandra X-ray Observatory Center, which is operated by the Smithsonian Astrophysical Observatory for and on behalf of the National Aeronautics Space Administration (NASA) under contract NAS8-03060. J.J.R., M.N., and D.H. acknowledge support from a Natural Sciences and Engineering Research Council of Canada (NSERC) Discovery Grant and a Fonds de recherche du Québec–Nature et Technologies (FRQNT) Nouveaux Chercheurs Grant. J.J.R. and M.N. acknowledge funding from the McGill Trottier Chair in Astrophysics and Cosmology. D.H. acknowledges support from the Canadian Institute for Advanced Research (CIFAR). P.A.E. acknowledges UKSA support.

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References

Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017a, PhRvL, 119, 161101
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017b, ApJL, 848, L12
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017c, ApJL, 848, L13
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017d, Natur, 551, 85
Alexander, K. D., Berger, E., Fong, W., et al. 2017, ApJL, 848, L21
Arcavi, I., Hosseinzadeh, G., Howell, D. A., et al. 2017, Natur, 551, 64
Arnaud, K. A. 1996, in ASP Conf. Ser. 101, Astronomical Data Analysis Software and Systems V, ed. G. H. Jacoby & J. Barnes (San Francisco, CA: ASP), 17
Coulter, D. A., Foley, R. I., Kilpatrick, C. D., et al. 2017, Sci, 358, 1556
Cowperthwaite, P. S., Berger, E., Villar, V. A., et al. 2017, ApJL, 848, L17
Drout, M. R., Piro, A. L., Shappee, B. J., et al. 2017, Sci, https://doi.org/10.1126/science.aao0049
Evans, P. A., Cenko, S. B., Kennea, J. A., et al. 2017, Sci, https://doi.org/10.1126/science.aap9580
Fong, W., Berger, E., Blanchard, P. K., et al. 2017, ApJL, 848, L23
Fruscione, A., McDowell, J. C., Allen, G. E., et al. 2006, Proc. SPIE, 6270, 62701V
Goldstein, A., Veres, P., Burns, E., et al. 2017, ApJL, 848, L14
Gottlieb, O., Nakar, E., & Piran, T. 2018, MNRAS, 473, 576
Gottlieb, O., Nakar, E., Piran, T., & Hotokezaka, K. 2017, arXiv:1710.05896
Guidorzi, C., Margutti, R., Brout, D., et al. 2017, ApJL, 851, L36
Haggard, D., Nynka, M., Ruan, J. J., et al. 2017, ApJL, 848, L25
Hallinan, G., Corsi, A., Mooley, K. P., et al. 2017, Sci, https://doi.org/10.1126/science.aap9585
Hjorth, J., Levan, A. J., Tanvir, N. R., et al. 2017, ApJL, 848, L31
Im, M., Yoon, Y., Lee, S.-K. J., et al. 2017, ApJL, 849, L16
Kasliwal, M. M., Nakar, E., Singer, L. P., et al. 2017, Sci, https://doi.org/10.1126/science.aap9455
Lazzati, D., Deich, A., Morsony, B. J., & Workman, J. C. 2017a, MNRAS, 471, 1652
Lazzati, D., Perna, R., Morsony, B. J., et al. 2017b, arXiv:1712.03237
Levan, A. J., Lyman, J. D., Tanvir, N. R., et al. 2017, ApJL, 848, L28
Margutti, R., Berger, E., Fong, W., et al. 2017, ApJL, 848, L20
McCully, C., Hiramatsu, D., Howell, D. A., et al. 2017, ApJL, 848, L32
Mooley, K. P., Nakar, E., Hotokezaka, K., et al. 2017, arXiv:1711.11573
Murguia-Berthier, A., Ramirez-Ruiz, E., Kilpatrick, C. D., et al. 2017, ApJL, 848, L34
Nakar, E., & Piran, T. 2017, ApJL, 834, 28
Pian, E., D’Avanzo, P., Benetti, S., et al. 2017, Natur, 551, 67
Piro, A. L., & Kollmeier, J. A. 2017, arXiv:1710.05822
Savchenko, V., Ferrigno, C., Kuulkers, E., et al. 2017, ApJL, 848, L15
Schlaufly, E. F., & Finkbeiner, D. P. 2011, ApJ, 737, 103
Shappee, B. J., Simon, J. D., Drout, M. R., et al. 2017, Sci, in press (arXiv:1710.05432)
Smartt, S. J., Chen, T.-W., Jerkstrand, A., et al. 2017, Natur, 551, 75
Soares-Santos, M., Holz, D. E., Annis, J., et al. 2017, ApJL, 848, L16
Troja, E., Piro, L., van Eerten, H., et al. 2017, Natur, 551, 71
Valenti, S., Sand, D. J., Yang, S., et al. 2017, ApJL, 848, L24
Vernon, D. A., Ferland, G. J., Korista, K. T., & Yakovlev, D. G. 1996, ApJ, 465, 487
Wilms, J., Allen, A., & McCray, R. 2000, ApJ, 542, 914

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Erratum: “Brightening X-Ray Emission from GW170817/GRB 170817A: Further Evidence for an Outflow” (2018, ApJL, 853, L4)

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Received 2018 May 2; published 2018 May 23

1 Introduction

We correct an error in Ruan et al. (2018), wherein the values for the X-ray photon index $\Gamma$ and the absorbed X-ray flux $F_{0.3-8 \text{ keV}}$ are inadvertently switched for GW170817/GRB 170817A at 109.2 days post-merger. This typo appears in Table 1 and in the text of the published article. Below, we provide an updated version of Table 1 containing the correct values. The conclusions of the article are not affected. We thank Harvey Tananbaum for pointing out this error.

Table 1

| Source ID | Days Post-burst | Count Rate (0.5–10 keV) $(10^{-4} \text{ cts s}^{-1})$ | Power-law Index $\Gamma$ | Flux (0.3–8 keV) $(10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2})$ | Luminosity (0.3–10 keV) $(10^{38} \text{ erg s}^{-1})$ | $\chi^2$/dof |
|-----------|----------------|-------------------------------------------------|--------------------------|-------------------------------------------------|-------------------------------------------------|----------------|
| GRB170817A | 15.6 | 3.5 ± 0.7 | 2.4±0.8 | 0.36±0.07 | 10.4±2.6 | 3.15/5 |
| | 109.2 | 14.7 ± 1.3 | 1.5±0.23 | 1.6±0.27 | 42.5±3.7 | 11.9/13 |
| CXOU J130948 | 15.6 | 2.5 ± 0.6 | 1.3±0.8 | 0.4±0.09 | 10.8±2.4 | 1.77/3 |
| | 109.2 | 1.2 ± 0.4 | 2.4±1.7 | 0.7±0.1 | 6.5±9.9 | 0.26/2 |
| CXOU J130946 | 15.6 | 3.3 ± 0.7 | –0.4±0.8 | 1.1±0.1 | 38.8±9.8 | 3.68/5 |
| | 109.2 | 2.5 ± 0.7 | 0.4±0.9 | 0.5±0.5 | 17.5±11.9 | 2.41/3 |
| NGC 4993 | 15.6 | 11.0 ± 1.2 | 1.5 ± 0.4 | 1.3 ± 0.2 | 34.8±3.2 | 12.73/11 |
| | 109.2 | 12.3 ± 1.2 | 1.4±0.3 | 1.4±0.2 | 38.9±5.8 | 13.09/11 |

Note. All reported errors represent 90% confidence intervals. The neutral hydrogen absorption was frozen to $N_H = 7.5 \times 10^{20} \text{ cm}^{-2}$ for all spectral fits, based on NGC 4993’s $A_V = 0.338$ (Schlafly & Finkbeiner 2011). A luminosity distance of 42.5 Mpc was assumed for all sources.

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References

Ruan, J. J., Nynka, M., Haggard, D., Kalogera, V., & Evans, P. 2018, ApJL, 853, L4
Schlafly, E. F., & Finkbeiner, D. P. 2011, ApJ, 737, 103