Widespread irreversible changes in surface temperature and precipitation in response to CO$_2$ forcing

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Some climate variables do not show the same response to declining atmospheric CO$_2$ concentrations as before the preceding increase. A comprehensive understanding of this hysteresis effect and its regional patterns is, however, lacking. Here we use an Earth system model with an idealized CO$_2$ removal scenario to show that surface temperature and precipitation exhibit globally widespread irreversible changes over a timespan of centuries. To explore the climate hysteresis and reversibility on a regional scale, we develop a quantification method that visualizes their spatial patterns. Our experiments project that 89% and 58% of the global area experiences irreversible changes in surface temperature and precipitation, respectively. Strong irreversible response of surface temperature is found in the Southern Ocean, Arctic and North Atlantic Ocean and of precipitation in the tropical Pacific, global monsoon regions and the Himalayas. These global hotspots of irreversible changes can indicate elevated risks of negative impacts on developing countries.

In recent decades, a series of climate reversibility studies have been conducted including the Carbon Dioxide Removal Model Intercomparison Project$^5$. Hysteresis and irreversibility were identified in the Atlantic Meridional Overturning Circulation$^{3-5}$, Antarctica Ice Sheet$^{6,7}$, El Niño–Southern Oscillation$^8$, Intertropical Convergence Zone$^9$, East Asian monsoon$^{10}$, sea level$^{11}$ and global mean surface temperature$^{12-17}$ and precipitation$^{12,13,18}$. These examples raise the question of how hysteresis and reversibility emerge at a regional level beyond the individual climate components. Our perception of climate change is based on a regional scale$^{19}$. Therefore, understanding climate hysteresis and reversibility at a regional level is crucial for assessing climate recoverability from the rise of anthropogenic global warming and for designing climate policy$^{20}$. Despite their importance, the hysteresis and reversibility of regional climate have not yet been clearly addressed so far.

In this Article, we aim to address the following key questions: how do the global patterns of hysteresis and reversibility look; which specific regions of the world exhibit irreversible changes; what is the key mechanism that shapes global hysteresis and reversibility patterns? We specifically focus on the response of two key climatic variables—surface temperature and precipitation—to CO$_2$ forcing.

For our investigation, we conducted the CO$_2$ ramp-up and ramp-down experiment using the Community Earth System Model (CESM) with 28 ensemble members. In the experiment, the atmospheric CO$_2$ concentration level is increased from 367 ppm to 1,468 ppm (four times the initial level) at 1% per year, then decreased back to 367 ppm at the same rate. The details of the experiment can be found in Methods.

Framework to quantify hysteresis and reversibility

A suitable method is required to analyse spatial patterns of hysteresis and reversibility from the experiment results. The key problem is quantifying hysteresis and reversibility for an entire climate trajectory spanning the CO$_2$ ramp-up and -down scenario. Previous studies have relied on analysis of the difference between two specific periods with the same CO$_2$ level (for example, refs. $^{12,16,21}$). However, this method cannot completely measure the hysteresis and reversibility for an entire duration of the simulation spanning a wide range of CO$_2$ levels. Here we develop a conceptual framework to quantify climate hysteresis and reversibility.

A trajectory of the climate state (for example, surface temperature or precipitation) under the CO$_2$ forcing can be represented as a loop in the CO$_2$ phase space (Fig. 1a). Under the scenario, the CO$_2$ concentration ($F$) is increased from the present-day level ($F_{\text{presp}}$) to the future peak level ($F_{\text{peak}}$) and restored back to the present-day level ($F_{\text{presp}}$). If the system exhibits hysteresis, the climate trajectory...
During CO₂ increase (x\text{down}) and decrease (x\text{up}) would be different. Here hysteresis for the entire trajectory can be measured as the area of the loop, defined as the hysteresis area (A):

\[
A = \int_{F_{\text{present}}}^{F_{\text{peak}}} |x_{\text{down}}(F) - x_{\text{up}}(F)| \, dF. \tag{1}
\]

The absolute difference between x\text{down} and x\text{up} holds the hysteresis area always positive. Reversibility of a system can be measured as whether the trajectory return to its initial state, indicated as an open-loop (irreversible change) and closed-loop (reversible change):

\[
\text{Loop} = \begin{cases} 
\text{Open} & \text{for } x_{\text{down}}(F_{\text{present}}) \neq x_{\text{up}}(F_{\text{present}}) \\
\text{Closed} & \text{for } x_{\text{down}}(F_{\text{present}}) = x_{\text{up}}(F_{\text{present}})
\end{cases} \tag{2}
\]

The definitions of hysteresis and reversibility are schematically illustrated in Fig. 1a. These definitions can quantify hysteresis and reversibility as a single scalar value for the entire climate trajectory; therefore, using them, we can map the hysteresis area and openness of the loop for each grid point of the experiment result. The examples of the various types of loop characterized by the hysteresis area and openness are shown in Fig. 1b. If the system has no hysteresis at all, it would exhibit no hysteresis area (A = 0) and would be closed-loop. It should be noted that the illustrated examples show possible cases of trajectory in a conceptual level. In physical perspective, all types of loop may not be equally plausible if the hysteresis area and openness are correlated. The detailed physical linkage between them is explained in later parts of the paper.

Note that an open-loop trajectory does not always indicate that a system is completely irreversible. Even if the loop is open, there is a possibility that the system will returns to its initial state if sufficient time is provided after the forcing reaches the initial level\(^2\). Nevertheless, at least, they show that the climate system cannot be immediately restored to its initial state even after successful removal of the atmospheric CO₂. The soft definition of irreversibility provides a practical classification for climate recoverability within a human-perceptible timescale.

**Hysteresis maps for surface temperature and precipitation**

Using the framework, we explore spatial patterns of hysteresis and reversibility in response to CO₂ forcing. We calculated the hysteresis area and loop openness for each grid point of the simulated surface temperature and precipitation. For the loop openness, we applied the refined definition of equation (2), which additionally considers the intersection of loop. The resultant hysteresis map depicts the hysteresis area as colour shading and hatched regions of a closed-loop response (Figs. 2a and 3a). The hysteresis map for the surface temperature is presented in Fig. 2a. The hysteresis area of the surface temperature shows large regional differences over the globe. It is the largest in the Southern Ocean near the coast of Antarctica (A = 5,440 °C ppm).
and smallest in the central Atlantic Ocean near the Bermuda islands ($A = 128$ °C ppm). Notably, the second-largest hysteresis emerges in the North Atlantic Ocean. The Northern and Southern hemispheres show contrasting responses to CO$_2$ forcing; the Southern Hemisphere shows larger hysteresis than the Northern Hemisphere (Fig. 2b). The ocean shows a similar hysteresis area to land, but it is exceptionally larger in the North Atlantic Ocean and Southern Ocean (Fig. 2b). Approximately 89% of the total global area show an open-loop response to CO$_2$ forcing (Fig. 2a). Only specific regions in the northern mid-latitudes and the Arctic show a closed-loop response (~11% of the total global area). This indicates that a vast majority of regions over the globe experience irreversible changes in surface temperature in response to CO$_2$ forcing. A strong correlation between hysteresis and reversibility is found; the closed-loop (open-loop) regions tend to have a small (large) hysteresis area (Supplementary Fig. 1).

Precipitation shows large regional differences in hysteresis and reversibility (Fig. 3a). The hysteresis area is the largest over the equatorial tropical Pacific where precipitation is heavily influenced by the Intertropical Convergence Zone ($A = 3,080$ mm per day ppm). Large precipitation hysteresis also emerges in the North Atlantic Ocean, Himalayas and Indonesia. The western Sahara Desert shows the smallest hysteresis ($A = 2.5$ mm per day ppm). The ocean shows a larger globally averaged precipitation hysteresis than land (Fig. 3b). The Northern and Southern hemispheres show a quasi-symmetric distribution in the zonal mean hysteresis area except in the North Atlantic Ocean (Fig. 3b). Approximately 58% of the global area shows an open-loop response for precipitation, and the rest of the 42% shows a closed-loop response (Fig. 3a). Thus, roughly half of the global area experiences irreversible changes in precipitation to CO$_2$ forcing, and the other half is reversible. The closed-loop response emerges in the Arctic, northern mid-latitude regions and Australia, and the open-loop response emerges in the equatorial tropical Pacific, Southern Ocean and Antarctica (Fig. 3a). Unlike the surface temperature, there is no distinct correlation between the hysteresis area and the openness of the loop (Supplementary Fig. 1).
To consider the seasonality of surface temperature and precipitation, we additionally calculated a hysteresis map for boreal summer and boreal winter (Supplementary Figs. 2–5). Compared with the annual pattern, the boreal summer and winter patterns for surface temperature hysteresis are different. The hysteresis area is increased in North Africa and India for boreal summer (Supplementary Fig. 2); and the Amazon rainforest and islands in the Arctic for boreal winter (Supplementary Fig. 3). In particular, the Arctic shows strong seasonal contrast in hysteresis and reversibility. The seasonal precipitation patterns also differ from the annual mean. In particular, it highlights irreversible response in global monsoon regions. A large hysteresis and open-loop response of precipitation emerges in India, Southeast Asia, East Asia and Northern Africa for boreal summer (Supplementary Fig. 4) and in western and eastern United States and western South America, southern Africa and Southeast Asia for boreal winter (Supplementary Fig. 5).

**Hotspots of irreversible changes**

Focusing on comprehensive response in climate variables, we identified hotspots of irreversible climate changes over the land based on the quantified hysteresis and reversibility. The hotspot for irreversible climate change is defined as a land region where both surface temperature and precipitation show an open-loop response with a large hysteresis area (Methods provide detailed criteria).

The hotspot map is presented in Fig. 4. Under our idealized CO2 forcing scenario, the hotspots emerge in most developing countries located in South America and Africa, such as Chile, Venezuela, Brazil, Nigeria, Ethiopia and Zimbabwe. Developing countries in Central America and South Asia are also classified as the hotspot. In contrast, only a few developed countries are classified as the hotspots, such as Ireland and New Zealand. The contrast between developing and developed countries implies a strong regional inequality in climate reversibility to anthropogenic global warming. The strong irreversible changes in developing countries translate into the long-lasting impact of anthropogenic CO2 emissions, which would vastly increase the risks of climate change with great social costs.

The hotspots also emerge in the regions covered with ice sheets, coastal regions of Antarctica, Greenland and Alaska. Large hysteresis and irreversible changes in these regions can potentially
cause the hysteresis in ice sheets2,24,29 and marine ecosystems30,31. The polar hotspots highlight the irreversibility of polar regions to CO₂ emissions.

**Mechanism of hysteresis and reversibility**

The detailed mechanisms for hysteresis and reversibility of surface temperature and precipitation are regionally dependent. For example, the large surface temperature hysteresis in the Southern Ocean is possibly related to warming delay effect mainly caused by the strong Ekman upwelling of the cold deep ocean water and heat uptake by the deep ocean23–25. The large precipitation hysteresis over the tropical Pacific can be explained by the hysteresis of the ITCZ due to the slow ocean circulation adjustment related to cross-equatorial energy exchanges between the Northern and Southern hemispheres.

Although the detailed dynamical process differs between regions, a fundamental mechanism for hysteresis and reversibility can be broadly modelled as an inertia effect. The multi-stability is also a possible mechanism for them26–28, but the inertia effect is a more reasonable hypothesis in this case (Supplementary Discussion 1).

The inertia induces a delayed response to CO₂ forcing and can result in hysteresis and irreversible changes. Here we present a simple inertia model for hysteresis and reversibility

\[ \frac{dx}{dt} = -\lambda (x - x_{eq}) \]  

\[ x_{eq} = \alpha F + \beta. \]

The climate state \( x \) is restored to equilibrium state \( x_{eq} \) at the rate of the inertia parameter \( \lambda \). The equilibrium state \( x_{eq} \) depends linearly on the forcing \( F \) with a sensitivity, \( \alpha \). \( \beta \) is an adjustment parameter for the \( x_{eq} \). \( \lambda \) characterizes the inertia of the system. A small (large) \( \lambda \) induces slow (fast) damping of the \( x \) to \( x_{eq} \), resulting in high (low) inertia of the system.

A solution of the simple inertia model is presented in Fig. 5a. The forcing \( F \) is increased from 367 to 1,468 with 1% per unit model time and decreased back to the initial level with the same rate as a parallel to the Earth system model experiment (Methods provide details). The parameter is set as \( \lambda = 0.02, \alpha = 1 \) and \( \beta = 1 \). The solution follows a different trajectory during increasing (orange line in Fig. 5a) and decreasing forcing (blue line in Fig. 5a) and exhibits hysteresis. Also, it does not return to its initial state, showing a difference between final and initial state, \( \Delta x \). The inertia effect characterized by \( \lambda \) induces such hysteresis and irreversible changes in response to the transient forcing. To examine the effect of \( \lambda \), the hysteresis area and \( \Delta x \) are calculated with varying \( \lambda \). As \( \lambda \) increases (or the inertia decreases), both hysteresis area and \( \Delta x \) decreased, converging into zero (black line in Fig. 5b,c). If \( \Delta x \) is sufficiently close to zero (\( \Delta x \approx 0 \)), the response becomes virtually reversible. This theoretical inertia curve provides the baseline to examine the inertia effect on hysteresis area and \( \Delta x \). A more detailed mathematical discussion is provided in Supplementary Discussion 2.

To examine how well the inertia effect explains the global hysteresis pattern, we compared the Earth system model experiment result with the theoretical inertia curve. For each grid point of the simulated surface temperature and precipitation, the \( \lambda \) is fitted (Supplementary Fig. 6), and normalized hysteresis area and \( \Delta x \) are calculated (Supplementary Figs. 7 and 8). The normalization excludes the effect of \( \alpha \) on the hysteresis area and \( \Delta x \), thus it allows us to directly examine the effect of \( \lambda \) (details can be found in Methods).

The hysteresis area of the surface temperature strongly follows the theoretical inertia curve \((R^2 = 0.84)\) (Fig. 5b). The \( \Delta x \) also follows it \((R^2 = 0.65)\) with some differences (Fig. 5c). There are two distinct branches on both hysteresis areas and \( \Delta x \). Most of the globe regions are on the main branch that follows the theoretical inertia curves well. As an exception, the North Atlantic Ocean appears on the secondary branch that does not follow the theoretical inertia curve. In the North Atlantic Ocean, the surface temperature response cannot be explained by the simple inertia model; it gradually warms during the CO₂ ramp up, followed by abrupt cooling and then warming during CO₂ ramp down (Fig. 2c). Physically, the temperature response is associated with the overshooting Atlantic Meridional Overturning Circulation during the CO₂ removal phase29,40.

The hysteresis area of the precipitation weakly follows the theoretical inertia curve \((R^2 = 0.39)\) (Fig. 5d), and the \( \Delta x \) does not follow...
Fig. 5 | The simple inertia model for hysteresis and reversibility. a, Solution of the simple inertia model. The x and F correspond to climate state (for example, surface temperature and precipitation) and CO$_2$ forcing, respectively. They are normalized to have a range of 0 to 1. The trajectory shows the climate trajectory during increasing F (orange line) and decreasing F (blue line). The black dashed line is a static equilibrium solution where F is fixed. b-e, Comparison of the theoretical inertia curve with the experiment result. The normalized hysteresis area and Δx are calculated with varying λ using the simple inertia model (thick black line). For each grid point of the experiment output, the x is fitted and normalized hysteresis area and Δx are calculated (Methods). The dots represent the result for each grid point and the colour shows the latitude of the dot. The results show normalized hysteresis area of surface temperature (b) and precipitation (d). The result for normalized Δx for surface temperature (c) and precipitation (e).

Discussion

The widespread hysteresis over the globe may induce hysteresis in diverse parts of the climate system. For example, the climate hysteresis in the Himalayas and Arctic can affect mountain glaciers$^{42}$ and summer sea ice$^{43-45}$, respectively, and the hysteresis in the Amazon can influence the rainforest$^{46}$.

The results presented in this paper are based on a forcing scenario that varies CO$_2$ concentration levels at a rate of 1% per year following the standard protocol of the Carbon Dioxide Removal Model Intercomparison Project$^2$ (Supplementary Fig. 9). The regional patterns of hysteresis and reversibility would depend on the rate of forcing (Supplementary Figs. 10 and 11). So far, the additional rate-sensitivity experiment showed that the surface temperature hysteresis responses over the Southern Ocean, North Atlantic and Arctic vary with the forcing rate (Supplementary Fig. 10), while the precipitation hysteresis is less sensitive to the rate of forcing (Supplementary Fig. 11). The dependency of regional climate hysteresis on the forcing rate needs to be further explored in the future.

Our quantification method can be universally applicable for any type of climate variable and forcing scenario, facilitating an understanding of regional patterns of hysteresis and reversibility. Further, the original definitions presented in this paper can be extended or modified; for example, the reversibility can be measured as Δx instead of openness of the loop to quantify the reversibility with more detail.

Our results show that the influence of greenhouse gas is not confined to the current warming period, and it extends beyond the human-perceptible timescale. Even after it is completely restored to the pre-industrial level, the climate may not be reversible in most regions, especially over many developing countries.

Online content

Any methods, additional references, Nature Research reporting summaries, source data, extended data, supplementary information, acknowledgements, peer review information; details of author contributions and competing interests; and statements of data and code availability are available at https://doi.org/10.1038/s41558-022-01452-z.

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Methods

Experiment configuration and design. We conducted the CO$_2$ ramp-up and ramp-down experiment using the Community Earth System Model version 1.2 (ref. 1) (CESM1.2). The model is composed of the atmosphere (Community Atmospheric Model version 5), ocean (Parallel Ocean Program version 2), sea ice (Community Ice Code version 4) and land models (Community Land Model version 4). The atmospheric model has a horizontal resolution of approximately 1° and 30 vertical levels. The ocean model has 60 vertical levels, with a longitudinal resolution of 1° and a gradually changing latitudinal resolution of 1/3° near the equator to 1/2° near the poles. The land model includes the carbon–nitrogen cycle.

The experiment was performed with idealized CO$_2$ scenarios composed of two different phases (Supplementary Fig. 9). In the first phase, the atmospheric CO$_2$ concentration level is fixed at 367 ppm as a present-day level, and the experiment is run for 900 years with a single ensemble (PD period). Before the PD period, the model is spun up for 800 years. In the second phase, the CO$_2$ concentration level increased from 367 ppm to 1,478 ppm at a rate of 1% per year for 140 years (ramp-up period), then decreased back to 367 ppm at the same rate for 140 years (ramp-down period). Followed by the ramp-down period, the CO$_2$ level is fixed at 367 ppm for 220 years (restoring period). The second phase experiment is run with 28 ensemble members with different initial conditions, which are extracted from the PD period. The initial conditions are set to have different phases of the multi-decadal climate oscillations such as the Atlantic Multidecadal Oscillation and Pacific Decadal Oscillation.

Calculation of the hysteresis area and openness. We calculated the hysteresis area and openness of the loop for each grid point of surface temperature and precipitation. We used ensemble mean surface temperature and precipitation. The ensemble mean can filter out the internal variability and show forced response to CO$_2$ forcing. The calculation covers CO$_2$ ramp-up and ramp-down periods. To remove high-frequency fluctuation, we applied a moving mean filter of 21 years.

The hysteresis area was calculated by equation (1) using the trapezoidal numerical integration method. The openness of the loop was calculated by a different but stricter definition than the one introduced in the main text (equation (2)). In the main text, for simplicity, we implicitly assumed that the trajectories during CO$_2$ ramp up ($x_{up}$) and ramp down ($x_{down}$) do not intersect each other, and they can only meet at the lowest level of CO$_2$ ($F_{lowest}$) only if the loop is closed (Fig. 1). However, in the several grid points of experimental output, the trajectories intersect at the intermediate level of CO$_2$ and exhibit a twisted curve shape. Such a type of loop is reasonable to be read as a closed-loop (reversible response) although it is characterized as open-loop by the definition in the main text (equation (2)). To account for these cases, we calculated the openness of the loop using a refined definition as follows:

\[
\text{Loop} = \begin{cases} 
\text{Open for min}(x_{up} - x_{down}) \times \max(x_{up} - x_{lowest}) > 0 & \text{if } x_{up} - x_{lowest} > 0 \\
\text{Closed for min}(x_{up} - x_{lowest}) \times \max(x_{up} - x_{down}) \leq 0 & \text{if } x_{up} - x_{lowest} \leq 0
\end{cases}
\]

(4)

where \(x_{up}\) and \(x_{down}\) are ramp-up and ramp-down trajectories during a range of CO$_2$ levels from \(F_{lowest}\) to \(F_{lowest}\), respectively. This definition classifies a trajectory as a closed-loop (open-loop) if \(x_{up} - x_{lowest}\) and \(x_{up} - x_{down}\) (do not) intersect during specific ranges of CO$_2$ levels from \(F_{lowest}\) to \(F_{lowest}\). The \(F_{lowest}\) is 367 ppm by the CO$_2$ scenarios in our experiment, and we selected \(F_{lowest}\) as 1,300 ppm. We directly applied the refined definition (equation (4)) to the processed experiment output.

Detecting hotspots of irreversible changes. The hotspot of irreversible changes is detected based on the annual mean land surface temperature and precipitation. The hotspot is defined as the land area where both surface temperature and precipitation show open-loop response and hysteresis area larger than its 50th percentile.

Solving the simple inertia model. The simple inertia model (equations (3a) and (3b)) was solved using the Euler time integration method with a time interval, \(\Delta t = 0.01\). The forcing, \(F(t)\) was imposed as:

\[
F(t) = \begin{cases} 
F_{up} & 0 \leq t < T/2 \\
F_{up} + T/2 & T/2 \leq t < T
\end{cases}
\]

(5)

where \(F_{up} = 367, r = 1.01, T = 280\) as a parallel to the Earth system model experiment. The model parameters are set as \(r = 0.22, \alpha = 1\) and \(\beta = 0\). The initial condition of \(x\) was set to \(F_{up}\). All quantities are dimensionless. To obtain the theoretical baseline of the inertia effect (black lines in Fig. 5b–e), we calculated the normalized hysteresis area and \(\Delta x\) with varying \(\lambda\). The normalized hysteresis area and \(\Delta x\) were calculated using the normalized x and F that have a range of 0 to 1.

Linking the simple inertia model to the experiment result. The hysteresis area and \(\Delta x\) of the surface temperature and precipitation were compared with the theoretical baseline obtained from the simple inertia model (black lines in Fig. 5b–e). The \(\lambda\) was fitted for each grid point of the surface temperature and precipitation (Supplementary Fig. 6). We used normalized surface temperature, precipitation and CO$_2$ concentration levels that have a range of 0 to 1. Substitute equation (3b) to equation (3a) and discretize \(\Delta t\) term with forward Euler scheme:

\[
\frac{X_{i+1} - X_i}{\Delta t} = -\lambda x_i + \lambda F_i + \lambda \beta
\]

(6)

where \(\Delta t\) is year and subscript \(i\) indicates ith time-step point (\(i = 1, 2, \ldots, 279\)). We performed linear regression for equation (6) and estimated \(\lambda\).

The normalized hysteresis area and \(\Delta x\) were calculated using the normalized x and F (Supplementary Figs. 7 and 8). The normalization can exclude the effect of \(\alpha\) on hysteresis area and \(\Delta x\); thus, it allows us to directly compare the Earth system model experiment result with the theoretical inertia baseline.

Specifically, \(\alpha\) linearly scales the hysteresis area and \(\Delta x\). If \(\alpha\) increases \(k\) times, both hysteresis area and \(\Delta x\) increase \(k\) times as well, \(\beta\) does not affect the hysteresis area and \(\Delta x\) at all; thus, we don’t have to consider it (Supplementary Discussion 2). Thus, the effect of \(\alpha\) on hysteresis area and \(\Delta x\) should be excluded to solely examine the effect of \(\lambda\). To eliminate its effect, for each grid point, we performed the linear regression that linearly transforms x and F to have a range of 0 to 1. It is essentially equivalent to setting \(\alpha = 1\) for all grid points; thus, the normalization excludes the effect of \(\alpha\) on the hysteresis area and \(\Delta x\).

Data availability

The data used in this paper are available from https://doi.org/10.6084/m9.figshare.2089123.v2.

Code availability

The code for the simple inertia model simulation is available from https://doi.org/10.6084/m9.figshear.20290017.v4.

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Author contributions

S.-K.K. conceived the idea of the research and drafted the article. S.-K.K. developed the conceptual framework for hysteresis and reversibility inspired by group discussion with J.S., H.-J.K., and N.I. J.S. conducted the Earth system model experiment. S.-K.K. and J.S. performed the post-data processing of the experiment output. S.-K.K. established the simple inertia model and performed the analysis on it. J.S. assisted the analysis. S.-K.K. prepared the figures by discussing with J.S. and S.-I.A. N.I. assisted in the illustration for Fig. 1. J.S. assisted in figure configuration for Fig. 2 and Fig. 3. S.-I.A. oversaw the project. All the authors discussed the results and wrote the article.

Competing interests

The authors declare no competing interests.

Additional information

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