1. INTRODUCTION

Characteristic timescales of variability provide an important way to probe the sub-parsec scale central engines in active galactic nuclei (AGNs) by providing information about the sizes and locations of emission regions. In blazars, i.e., BL Lacertae objects (BL Lacs) and flat spectrum radio quasars, Doppler-boosted emission from a relativistic jet has long been recognized to provide the only feasible explanation for their non-thermal spectra and radio morphologies on small scales (e.g., Blandford & Rees 1978; Urry & Padovani 1995). Still, quasi-periodic oscillations that have been claimed to be observed in any blazar at any wavelength. While these data are insufficient to strongly constrain models for such fluctuations, the presence of such a short timescale when the source is not in a very low state seems to favor the action of turbulence behind a shock in the blazar’s relativistic jet.

Key words: BL Lacertae objects: individual (S5 0716+714) – galaxies: active – galaxies: photometry

2. OBSERVATIONS AND DATA REDUCTION

Our observations of S5 0716+714 were carried out with an Andor EMCCD (Electron Multiplying Charge Coupled Device) camera mounted at the f/13 Cassegrain focus of the 1.2 m telescope operated by the Physical Research Laboratory (PRL) at Gurushikhar, Mt. Abu, India. We observed this source on 2008 December 23, 27, and 28 and 2009 January 3; the total amount of data collected over those four nights was 9.6 hr. The 1k × 1k EMCCD has square pixels with sides of 13 μm size. With electron multiplication technology, the read noise in the system is expected to be negligible compared with normal CCD cameras (Mackay et al. 2001) and the performance approaches near

QUASI-PERIODIC OSCILLATIONS OF \( \sim 15 \) MINUTES IN THE OPTICAL LIGHT CURVE OF THE BL LAC S5 0716+714

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ABSTRACT

Over the course of 3 hr on 2008 December 27, we obtained optical (\( R \) band) observations of the blazar S5 0716+714 at a very fast cadence of 10 s. Using several different techniques, we find fluctuations with an approximately 15 minute quasi-period to be present in the first portion of these data at a >3σ confidence level. This is the fastest quasi-periodic oscillation that has been claimed to be observed in any blazar at any wavelength. While these data are insufficient to strongly constrain models for such fluctuations, the presence of such a short timescale when the source is not in a very low state seems to favor the action of turbulence behind a shock in the blazar’s relativistic jet.

Key words: BL Lacertae objects: individual (S5 0716+714) – galaxies: active – galaxies: photometry

The bright, high-declination BL Lac, S5 0716+714, at redshift \( z = 0.31 \pm 0.08 \) (Nilsson et al. 2008) has been extensively studied across the electromagnetic spectrum and exhibits strong variability on a wide range of timescales, ranging from minutes to years (e.g., Gupta et al. 2008a, 2008b, 2009, and references therein). The optical duty cycle of S5 0716+714 is nearly unity, indicating that the source is always in an active state in the visible (Wagner & Witzel 1995). This blazar was recently shown to be a strong source in the high-energy gamma-ray band by Fermi (Abdo et al. 2009). Rapid fluctuations have long been known to characterize blazars, with the prototype, BL Lac, seen to flicker over just a few minutes in early single channel photometry with 15 s temporal resolution (Racine 1970).

Among the other blazars, PKS 2155−304 possibly showed a quasi-periodicity around 0.7 days during 5 days of observations at UV and optical wavelengths (Urry et al. 1993). Very recently, somewhat better evidence for a QPO of \( \sim 4.6 \) hr in the \( XMM-Newton \) X-ray light curve (LC) of PKS 2155−304 has been reported by Lachowicz et al. (2009). An \( XMM-Newton \) LC of the quasar 3C 273 appears to have a quasi-periodic component with a timescale of about 3.3 ks (Espaillat et al. 2008). Using the \( \sim 13 \) year long data taken by the All-Sky Monitor on the \( Rossi X-ray Timing \) \( X \)-ray satellite, Rani et al. (2009) reported good evidence of nearly periodic variations of \( \sim 17.7 \) days in the blazar AO 0235+164 and \( \sim 420 \) days in the blazar 1ES 2321+419. The narrow-line Seyfert 1 galaxy, RE J1034+396, while not a blazar, strongly indicated the presence of a \( \sim 1 \) hr periodicity during a 91 ks observation by the X-ray satellite \( XMM-Newton \) (Gierliński et al. 2008).

In this Letter, we exhibit evidence for a QPO of \( \sim 15 \) minutes in a single densely sampled optical LC of the blazar S5 0716+714. We first used a structure function (SF) analysis to find a hint of such a QPO and then quantified the strength of this signal using Lomb–Scargle Periodogram (LSP) and power spectral density (PSD) methods. We find this to be a strong case for the discovery of the shortest nearly periodic variation seen for any blazar, or for that matter, any AGN, in any wave band.
Figure 1. (a) R passband LC of the blazar S5 0716+714 in the first 1.66 hr of observations; the temporal origin is at 18.86976 hr UT on 2008 December 27. (b) The calibrated LC of the source for the entire observation, along with the differential instrumental magnitudes of standard stars 8 and 11, offset by 14.14 mag.

3. ANALYSIS AND RESULTS

In order to be certain that the apparent variability of S5 0716+714 is significant, we used the $F$-test, shown by de Diego (2010) to be superior to commonly used methods. The $F$-statistic is the ratio of the sample variances, or $F = \frac{s_Q^2}{s_S^2}$, where the variance for the quasar differential LC is $s_Q^2$, while that for the standard star is $s_S^2$. We used the $F$-test code available in R6 and find $F = 18.3748$, with a significance level of 0.9999998 or $>5\sigma$.

We have also calculated the variability amplitude parameter, $A$ (Heidt & Wagner 1996), to see the percentage variation in the LC of source. For S5 0716+714, we find $A = 16.9\%$. The calculated fractional rms variability amplitude for the LC (Vaughan et al. 2003) is $F_{\text{var}} = 15.45$.

A visual inspection of the LC for the first 2 hr shown in Figure 1(a) indicates a possible periodic modulation of the variability at about 900 s, along with a hint of even faster modulations at the very beginning of the observation. The calibrated LC for the entire 3 hr of measurements taken at a 10 s cadence, along with the differential instrumental magnitudes of standard stars 8 and 11 is displayed in Figure 1(b). The LC averaged over 30 s intervals folded at a putative period of 900 s is shown in Figure 2(a).

3.1. Structure Function

The first-order SF is a simple way to search for periodicities and timescales of variability in time series data trains (e.g., Simonetti et al. 1985). Here, we give only a very brief summary of the method (for details, refer to Rani et al. 2009). The first-order SF for a data train, $a$, is defined as

$$D_a^1(k) = \frac{1}{N_a^1(k)} \sum_{i=1}^{N} w(i) w(i + k) [a(i + k) - a(i)]^2,$$  

where $w(i)$ is the weight assigned to each data point, $a(i)$ is the $i$th data point, $N_a^1(k)$ is the number of valid data points within a lag $k$, and $N$ is the total number of data points. The $F$-test and $A$ parameter provide statistical significance and amplitude of the variability, while the structure function helps identify the periodicities and timescales of the variability.
where $k$ is the time lag, $N_1^k(k) = \sum w(i)w(i + k)$, and the weighting factor, $w(i)$, is 1, if a measurement exists for the $i$th interval and 0 otherwise. For a time series containing a periodic pattern, the SF curve shows minima at time lags equal to the period and its subharmonics (e.g., Lachowicz et al. 2006), although dips and wiggles in SFs are not always reliable indicators of timescales (Emmanoulopoulos et al. 2010). The SF analysis curve of the whole data set is displayed in Figure 2(c). The first dip and the seven cycles of its subsequent subharmonics correspond to a possible period of $927 \pm 30\text{ s}$.

### 3.2. Lomb—Scargle Periodogram

The LSP, introduced by Lomb (1976) and extended later by Scargle (1982), is an excellent technique for searching time series, as long as white noise, $P_N(f) \propto f^0$, is the dominant noise process. Press & Rybicki (1989) provided a more practical mathematical formulation. For the details of method and formulae see Rani et al. (2009) and references therein.

We used an online available R-language code for the LSP.\(^7\) The LSP analysis of the whole data set is displayed in Figure 2(b). The LSP analysis revealed the detection of significant frequency corresponding to a period of 904 s with a significance level of 0.999999977. Two questions usually arise concerning the validity of a periodogram result (Scargle 1982); the first is statistical and the second is spectral leakage. The statistical difficulty is mitigated by the good S/N of $\sim 35$ in our case. Spectral leakage, which is also known as aliasing, involves the spreading of periodogram power to other frequencies that are actually not present in the data. Since our data are uniformly sampled, there might be chances of aliasing. But as essentially the same period is confirmed by SF and PSD analyses the strong signal is very unlikely to arise in this fashion.

However, as LCs of most AGNs contain red noise as well as white noise, a more robust test is required to quantify the presence of a QPO.

### 3.3. Power Spectral Density

The PSD is a powerful tool to search for periodic signals in time series, including those contaminated by white noise and/or red noise (e.g., Vaughan 2005). We employed a PSD analysis method (Vaughan et al. 2003; Vaughan 2005) that is suitable for these types of LCs. First, as shown in Figure 3, we fit a single power law (SPL) to the calculated PSD, assuming it to have a form $P(f) \propto f^\alpha$ at low frequencies and then examined the significance of the frequency peak using the method of Vaughan (2005). This analysis indicates the presence of a QPO signal with peak frequency $\approx 0.001077\text{ Hz (or period } \approx 928\text{ s)}$, with a $3.4\sigma$ significance level. The calculated significance is global, i.e., corrected for the number of frequencies tested. The range of frequencies used for calculating the global significance of the QPO is $0.0002 \leq f \leq 0.002$ which amounts to 28 frequency bins. This range excludes frequencies that are significantly dominated by white noise.

We next checked the statistical significance of this QPO using Monte Carlo simulations. We generated a series of $10^4$ simulated LCs following a given SPL having the same number of bins, mean, and variance as the observed LC (Timmer & Koenig 1995) using an IDL code available online.\(^8\) The PSD analysis resulting from the simulated LCs using an SPL with the index

\(^7\) [http://research.stowers-institute.org/efg/2005/LombScargle](http://research.stowers-institute.org/efg/2005/LombScargle)

\(^8\) [http://astro.uni-tuebingen.de/software/idl/aitlib/timing/timmerlc.html](http://astro.uni-tuebingen.de/software/idl/aitlib/timing/timmerlc.html)
from the best fit to our data is compatible with the results shown in Figure 2 and indicates an average significance of 3.2σ. We also considered the alternative null hypothesis of a broken power law (BPL). The best-fitting BPL indices are, respectively, +0.39 and −0.8 above and below the break frequency of ∼0.0011 Hz. The nominal statistical significance of the QPO frequency in this case is ∼3.1σ. Finally, we performed PSD analyses of simulated LCs generated from BPLs and calculated periodograms for each of them, finding that the periodic signal was still significant at 3σ. Hence, we conclude that the observed QPO at a frequency of ∼0.001077 Hz is statistically significant, irrespective of the assumed model of continuum power.

4. DISCUSSION AND CONCLUSIONS

This discovery of a nearly periodic signal of ∼900 s in the optical R passband LC of the blazar S5 0716+714 adds a unique new point to the variability studies of blazars at intra-night timescales. The presence of seven cycles with a >3σ significance level allows us to make a strong claim for the shortest optical QPO detected so far.

The simplest possible explanation for such a short period might be the flux arising from hot spots or some other non-axisymmetric phenomenon related to the orbital motions that are close to the innermost stable circular orbit around a supermassive black hole (SMBH; e.g., Zhang & Bao 1991). Adopting z = 0.31 for S5 0716+714 (Nilsson et al. 2008) means that a 900 s period at the inner edge of a corotating disk corresponds to a SMBH mass of 1.5 × 10^6 M⊙ for a non-rotating black hole (BH) and 9.6 × 10^6 M⊙ for a maximally rotating BH (Gupta et al. 2009). If the source arises somewhat further out in the accretion disk, then the BH mass would be even less than these modest values.

However, since blazar jets are pointing very close to the line of sight of the observer (e.g., Urry & Padovani 1995), the emerging flux, particularly in active phases, is dominated by emission from jets. Turbulence behind a shock propagating down a jet (e.g., Marscher et al. 1992) is a very plausible way to produce dominant eddies whose turnover times can yield short-lived, quasi-periodic fluctuations in emission at different wavelengths. Since Doppler boosting will greatly amplify the very weak intrinsic flux variations produced by small changes in the magnetic field or relativistic electron density, these intrinsically weak fluctuations can be raised to the level at which they can be detected (e.g., Qian et al. 1991). This same Doppler boosting reduces the timescale at which these fluctuations are observed compared to the timescale they possess in the emission frame. Although it is difficult to quantify these effects precisely, this mechanism does seem to provide an excellent way to understand the type of short-lived optical intra-night variability with periods of tens of minutes seen here.

It is also possible that QPOs originate from a relativistic shock propagating down a jet that possesses a helical structure, as can be induced by magnetohydrodynamical instabilities (Hardee & Rosen 1999) or even through precession. Indeed, in some cases where radio jets can be resolved transversely using Very Long Baseline Interferometry, edge-brightened and non-axisymmetric structures are seen (e.g., M87, Ly et al. 2007; Mkn 501, Piner et al. 2009). A relativistic shock propagating down such a perturbed jet will induce significantly increased emission at the locations where the shock intersects with a region of enhanced magnetic field and/or electron density corresponding to such a non-axisymmetric structure. Because Doppler boosting is a sensitive function of viewing angle, substantial changes in amplitude of jet emission can be seen by the observer (Camenzind & Krockenberger 1992; Gopal-Krishna & Wiita 1992).

Therefore, the observed periodic component in the optical LC of S5 0716+714 might be attributed to the intersections of a relativistic shock with successive twists of a non-axisymmetric jet structure, though they would have to be surprisingly tight to yield such a short period.

We have analyzed the optical R passband LC of the well-known BL Lac S5 0716+714 observed on 2008 December 27 with a 10 s cadence that provided the best time resolution so far obtained for a blazar. All different analyses indicate that this LC contains a periodic component to its fluctuations of about 15 minutes. Although this particular BL Lac showed earlier evidence of periodic variations in radio through X-ray wave bands ranging from tens of minutes to several years, our new data provide the shortest known quasi-period yet detected in a blazar.

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Figure 3. PSD of S5 0716+714. P(f) is the best-fitting SPL with index −0.58 ± 0.06 (the dot-dashed lines are the calculated uncertainty in the model); a 3σ confidence limit and the white-noise level are shown.
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