Chaos as a Source of Complexity and Diversity in Evolution

Kunihiko KANEKO

Department of Pure and Applied Sciences
University of Tokyo, Komaba, Meguro-ku, Tokyo 153, JAPAN

Abstract

The relevance of chaos to evolution is discussed in the context of the origin and maintenance of diversity and complexity. Evolution to the edge of chaos is demonstrated in an imitation game. As an origin of diversity, dynamic clustering of identical chaotic elements, globally coupled each to other, is briefly reviewed. The clustering is extended to nonlinear dynamics on hypercubic lattices, which enables us to construct a self-organizing genetic algorithm. A mechanism of maintenance of diversity, “homeochaos”, is given in an ecological system with interaction among many species. Homeochaos provides a dynamic stability sustained by high-dimensional weak chaos. A novel mechanism of cell differentiation is presented, based on dynamic clustering. Here, a new concept – “open chaos” – is proposed for the instability in a dynamical system with growing degrees of freedom. It is suggested that studies based on interacting chaotic elements can replace both top-down and bottom-up approaches.

1 Complexity, Diversity, and Emergence

Why are we interested in the effort to create “life-like” behavior in computers? The answers can be diverse, but my interest in such artificial biology lies in the construction of systems exhibiting the emergence and maintenance of complexity and diversity, in order to understand the evolution of the complex "society" of life. This problem is not so trivial, indeed. It is often difficult to conclude that a system’s emergent complexity is somewhat beyond that which would be expected on the basis of the rules explicitly implemented within a model [1]. Often, what people call “emergent” behavior comes from the lack of a full understanding of what is implied by the rules implemented in the model. In evolution, there is a stage of the emergence of novel features as well as a stage of slow-scale change of existing features. Gradual evolution after the emergence of a novel feature is often studied analytically with the use of stochastic differential equations, as, for example, is demonstrated by the neutral theory of evolution [2]. The “origin” of features, on the other hand, is often a difficult problem to solve analytically. The origins of life, eukaryotes, multi-cellular organism, germ-line segregation, and sex are examples of the emergence of such novel features. For such problems, we require a mechanism for how complex, higher-level behavior emerges from low-level interactions, without the implementation of explicit rules for such emergence. Such emergence, we believe, occurs through strong nonlinear interactions among the agents at the lower level. Nonlinear interaction among agents
often leads to chaotic behavior, which, we believe, can cause “aufheben” (German terminology of dialectic philosophy) to higher level dynamics, by which lower level conflicts are resolved. In this overview, we try to demonstrate that chaos is relevant to the emergence and maintenance of complexity and diversity. Chaos is the most universal mechanism to create complexity from simple rules and initial conditions. As will be seen, chaos can be a source of diversity: identical elements differentiate through chaotic dynamics. Through a dynamical process with instability, chaos also has the potentiality to create a higher-level dynamics. Problems we address in the present overview are as follows; (i) evolution to complexity, (ii) sources of diversity, (iii) maintenance of diversity, and (iv) successive creation of novelty and open-ended evolution to diversity. In §2, evolution to the edge of chaos, a complex state between chaos and order (a window), is studied with the use of an imitation game. An explicit example of the evolution of complexity is given. General concepts in globally coupled dynamical systems are briefly given in §3, including the dynamic clustering of synchronization, hidden coherence, and chaotic itinerancy. In §4 these novel concepts are applied to dynamical systems on a hypercubic lattice, which enables us to construct spontaneous genetic algorithms. A new concept – “homeochaos” – representing dynamical stability involving weak chaos with many degrees of freedom, is given in §5, as well as its relevance to various biological networks. Homeochaos (to be contrasted with ”homeostasis”) can provide for the maintenance of diversity. In §6, the concept of clustering in globally coupled maps is extended to the problem of cell division and differentiation. Here the novel concept of “open chaos” is demonstrated in a system with growing degrees of freedom. Open chaos leads to the formation of disparities in activities among cells, leading to the emergence of diversity and novelty. In §7 we show the advantage of our approach over top-down and bottom-up approaches.

2 Edge of Chaos in an Imitation game: Chaos as a source of Complexity

The increase of complexity through evolution is believed to be seen in many biological systems, not only in the hierarchical organization in genotypes and phenotypes but also in animal behavior and communication. A direction for the increase of complexity has recently been discussed as ”evolution to the edge of chaos” [2], since complexity is believed to be large at the border between order and chaos [3, 4, 5, 6, 7, 8]. However, there has been no clear simple example providing evidence of evolution to the edge of chaos, in the exact sense of dynamical systems theory. Chaos is defined only on dynamical systems with a continuous state, and is not defined for discrete-state systems such as cellular automata, which have been adopted for studies of the edge of chaos so far. Recently Suzuki and the author have presented an example of evolution to the edge of chaos by introducing a simple model for an imitation game of a bird song [9]. A bird song, for example, is known to increase its complexity through evolution and development (with more repertoire made up from combinations of simple phrases) [11]. A bird with a complex song is stronger in defending its territory [12]. Based on this observation of a function of bird song for the defense of territory, we have introduced an imitation game [10], in which the player who imitates the other’s song better wins the game. As a ”song”, a time series generated by a simple mapping $x_{n+1} = f(x_n) = 1 - ax_n^2$ (the logistic map) is adopted. As is well known the attractor of the map shows a bifurcation sequence from a fixed point, to cycles with periods 2, 4, 8, ..., and to chaos as the parameter $a$ is increased [13]. Here the parameter value $a$ is assumed to be different for each individual ”bird”. By this
choice of a song generator, one can examine whether a song evolves towards the edge of chaos. Each bird player $i$ chooses an initial condition, so that the time series of its own dynamics $x_{n+1}(i) = f_i(x_i) = 1 - a(i)x_n(i)^2$ can imitate the song of another player. For “preparation” of an initial condition, the player $i$ uses a feedback from the other player $j$’s song by $x_{n+1}(i) = f_i[(1 - \epsilon)x_n(i) + \epsilon x_n(j)]$ over a number of given time steps $t_{trs}$. Of course, birds have to choose initial conditions for starting the above feedback process, and also for singing, and the result of a game can depend on this choice of initial conditions [10]. Here we assume that they choose initial conditions randomly over $[-1, 1]$. Thus the game is probabilistic, although “strong” players (to be discussed) often win against “weak ones” with probability close (or equal) to one [10]. Repeating the imitation process, the distance between two songs $D(j, i) = \sum_{n=t_{trs}}^{t_{trs}+T}(x(j) - x(i))^2$ is measured. By changing the role of the player $i$ and $j$, $D(j, i)$ is measured. If $D(i, j) < D(j, i)$, the player $i$ imitates better than $j$, thus being the winner of the game, and vice versa. By reproducing the players according to their scores in the game, and by including mutation of the parameter $a$ [9], we have examined the dynamical states to which the songs evolve. Temporal evolution of the average of the parameter $a$ over all players shows successive plateaus, until it reaches $a \approx 1.94$, where it then remains. Plateaus corresponding to period-doubling bifurcation points or to the edge between periodic windows and chaos are observed successively. In Fig.1, the average score of players is plotted as a function of $a$. The score has a peak at the edge between chaos and windows for stable periodic cycles. To study evolution to the edge of chaos, we have also plotted the score of birds as a function of the Lyapunov exponent $\lambda$ for the dynamics $x \rightarrow f(x)$ (see Fig.2). Indeed, the score has a broad peak around $\lambda = 0$, corresponding to the edge of chaos. Thus, evolution of a song towards the edge of chaos is observed. The final value of $a \approx 1.94$ corresponds to the borderline between a periodic window (of period 4) and chaos. Besides evolution to the edge of chaos, it should be noted that the ”edge” reached by evolution lies between a periodic window and chaos. At a window, the dynamics show chaotic transients before attraction to a stable cycle. Here, a variety of unstable cycles coexist [13], which provide for a larger variety of dynamics, as transients. Transient chaos is important for the adaptation to a wide range of external dynamics. Evolution to the edge of windows may be a more robust and important concept than evolution to the edge of chaos. Escape from imitation could be a trigger for the evolution of complexity in many fields, such as in the evolution of Batesian mimicry [14], where one of the groups can survive better by imitating the pattern of another group, while the second group’s advantage in survival is lost if it is not distinguished well from the first. The increase of complexity of the patterns of some butterflies may be due to this “imitation” pressure. Another possible application is seen in the evolution of a communication code only within a given group. Studies of the evolution of such signals will be important in the future, from the viewpoint of complexity via chaotic dynamics.

3 Key Concept for the Origin of Complexity and Diversity: Dynamic Clustering in Networks of Chaotic Elements

To study the emergence of diversity, we need a mechanism by which identical elements differentiate into different groups spontaneously. Networks of chaotic elements, globally coupled to each other, provide an example of such a mechanism.
In many biological networks, the interaction among elements is not local but global. The simplest case of global interaction is studied as a “globally coupled map” (GCM) of chaotic elements. An example is given by

\[ x_{n+1}(i) = (1 - \epsilon)f(x_n(i)) + \frac{\epsilon}{N} \sum_{j=1}^{N} f(x_n(j)) \]  

where \( n \) is a discrete time step and \( i \) is the index of an element (\( i = 1, 2, \cdots, N = \) system size), and \( f(x) = 1 - ax^2 \) \(^{[15]}\). Without coupling (i.e., for \( \epsilon = 0 \)), each element shows chaotic behavior if \( a \) is large enough. The model is a mean-field-theory-type extension of coupled map lattices (CML). The above dynamics consists of parallel nonlinear transformation with a feedback from the “mean-field” \(^{[17]}\). In real biology, elements are not necessarily identical. The reason we start from identical elements is that we are interested in the origin of differentiation and diversity. That is, we are interested in the question: How can a set of identical units evolve to groups with different (dynamical) states? Through interaction, some elements oscillate synchronously, while chaotic instability gives a tendency for the destruction of coherence. Attractors in GCM are classified by the number of synchronized clusters \( k \) and the number of elements within each cluster \( N_k \). Here, a cluster is defined as the set of elements in which \( x(i) = x(j) \). Identical elements split into clusters with different frequencies, phases, or amplitudes of oscillation. Each attractor is coded by the clustering condition \([k, (N_1, N_2, \cdots, N_k)]\). In a globally coupled chaotic system in general, the following phases appear successively with the increase of nonlinearity in the system \(^{[14]}\): (i) **Coherent phase**: A coherent attractor (\( k = 1 \)) has occupied (almost) all basin volumes. (ii) **Ordered phase**: Attractors (\( k = o(N) \)) with few clusters have occupied (almost) all basin volumes. (iii) **Partially ordered phase**: Coexistence of attractors with many clusters (\( k = O(N) \)) and attractors with few clusters. (iv) **Turbulent phase**: All attractors have many clusters (\( k = O(N) \)); in most cases \( k \approx N \).

In the turbulent phase, although \( x(i) \) takes almost random values almost independently, there remains some coherence among elements. The distribution \( P(h) \) of the mean field \( h_n \equiv (1/N) \sum_j f(x(j)) \), sampled over long time steps, does not obey the law of large numbers \(^{[14]}\). The emergence of hidden coherence is a general property in a globally coupled chaotic system. This hidden coherence may be interesting in relation to EEG’s, where one measures a given average of neuronal (electric) activity. Although the firing of each neuron is not regular (i.e., chaotic or random), the amplitude of some average (EEG) still has a large enough amplitude of variation to be observed, which may suggest the existence of hidden coherence as in the above. In the partially ordered phase, orbits make itinerance over ordered states via highly chaotic states. In the ordered states the motion is partially coherent. Our system exhibits intermittent change between self-organization towards a coherent structure, and its collapse to a high-dimensional disordered motion. This dynamics, called **chaotic itinerancy**, has been found in a model of neural dynamics by Tsuda \(^{[18]}\), optical turbulence \(^{[19]}\), and in GCM \(^{[14]}\). Here, a number of ruins of low-dimensional attractors coexist in the phase space. The total dynamics consists of residencies at ruins interspersed with excursions into high-dimensional chaotic states. In the chaotic itinerancy in GCM, the degree of synchronization between two elements changes with time. Elements 1 and 2, for example may be almost synchronized for some time span, until desynchronization by high-dimensional chaos destroys the relationship. After some time, element 1 may be almost synchronized with element 5, for example, and so forth. Thus, the relationship between elements is dynamically changing. Indeed, such change of relationship is discussed in neural physiology \(^{[20]}\).
4 Clustering in hypercubic coupled maps; self-organizing genetic algorithms

Let us discuss an extension of the idea in §3 to population dynamics with mutation. The process of mutation is characterized by a diffusion process in the space of genes. If the “gene” space is represented by a bit space (such as \( i = 0010111 \), as is often the case for genetic algorithms [9]), the single point mutation process is given by a flip-flop \( 0 \leftrightarrow 1 \) at each position. Let us represent the population (density) of each species \( i \) by \( x_n(i) \), the mutation process is given by a diffusion in hypercubic bit space. When the population dynamics is represented by \( x(i) \rightarrow f(x(i)) \), the total dynamics is given by

\[
x_{n+1}(i) = (1 - \epsilon) f(x_n(i)) + \frac{\epsilon}{k} \sum_{j=1}^{k} f(x_n(\sigma_j(i))),
\]

where \( \sigma_j(i) \) is a species whose \( j \)’th bit is different from the species \( i \) (with only one bit difference), and \( k \) is the total bit length of species (total species is \( 2^k \) [22]). We use a binary representation to denote the lattice here; for example, site 42 for \( k = 6 \) means the hypercubic lattice point 101010. The above model is rather close to the model in the last section; instead of the global coupling in (1), nearest neighbor coupling on the hypercube is adopted here. In the model (2), we have again found the formation of synchronized clusters as in §3 (i.e., \( x_n(i) = x_n(j) \) for two elements \( i \) and \( j \) in the cluster). In the present case, the split to clusters is organized according to the hypercubic structure. Examples of such clusters follow. (a) 1 bit clustering Two clusters with synchronized oscillation are formed. Each of the clusters has \( N/2 = 2^{k-1} \) elements, determined by the bit structure. For example, elements may be grouped into two clusters with \(*0****\) and \(*1****\), (* means that the symbol there is either one or zero), each of which has \(2^{k-1}\) species. This clustering is formed by cutting the \( k \)-dimensional hypercube by a hyperplane. In the genetic algorithm [9], irrelevant bits are initially determined as “don’t care” bits represented by “#”. Here, such bits are spontaneously created with the temporal evolution. (b) 2 bit clustering Depending on initial conditions and parameters, both the number of clusters and the number of bits relevant to clustering can be larger than in case (a). An example is a 2-cluster state with 2 relevant bits by XOR (exclusive-or) construction. Here the elements split into the groups (i) **1*** or **0*** and (ii) **00*** or **11***, for example. (c) parity check clustering Elements split into two groups according to the parity of the number of 1’s in each bit representation. For example, elements split into two clusters as follows (i) 000, 011, 101, 110 and (ii) 001 010, 100, 111, for \( k = 3 \). The clustering, thus gives a parity check. It is a hypercubic version of the zigzag (1-dim) or checkerboard (2-dim) pattern [8]. Besides these examples, attractors with many clusters are also found. Most of these states are constructed by combining the above clustering schemes. For example, 4 clusters with two relevant bits are found as a direct product state of case (a). Here the hypercubic space is cut by two hyperplanes. Elements split into four clusters, for example, coded by 01****, 10****, 11****, and 00****. More complex examples are reported in [21]. Here we have to note that not all partitions are possible in the present case. Even if we start from an initial condition with an arbitrary clustering, the synchronization condition \( x(i) = x(j) \) for \( i, j \) belonging to the same cluster) is not satisfied at the next step for most such initial conditions. In contrast with the GCM case, not all possible partitions can be a (stable or unstable) solution of the evolution equation. As is discussed, the present result opens up the possibility of automatic genetic algorithms. Relevant bits are spontaneously formed. Furthermore we have found a chaotic
itinerancy state, where relevant bits change according to temporal evolution. In Fig. 3, the change of relevant bits for clustering is clearly seen. At stage A, two clusters are formed by the first bit (i.e., clusters 0\ldots\ldots\ldots\ldots and 1\ldots\ldots\ldots), and the second bit at stage B (\ldots\ldots\ldots\ldots and \ldots\ldots\ldots), and so on. With the introduction of external inputs to each element, it is also possible to have a clustered state following the external information [21]. Relevant information is extracted through this process spontaneously, which is stored as a relevant bit in the clustering. An application of the present clustering to “real life” will be found in the quasispecies of viruses [26]. As Eigen et al. discuss, viruses form quasispecies coded in hypercubic space. By taking account of population dynamics, the present clustering may give a theoretical basis for the dynamic and hierarchical grouping of quasispecies.

5 Maintenance of Diversity and Dynamic Stability: Homeochaos

As for the evolution to complexity, the notion of the “edge of chaos” in §2 is rather special. First, it is provided by a critical state, and should be sustained at a very narrow region (or at a single critical point) of the parameter space. Second, the system is given by a low-dimensional dynamical system, i.e., with very few degrees of freedom. Hence the notion “edge of chaos” is insufficient to understand the diversity and complexity of a biological system. In an ecological system, many species are under strong nonlinear interaction, and keep some kind of stability with diversity. This is not easily sustained. We also have to mention that static equilibria with many species are usually unstable, as studied by May [25] in a random network model. Thus it is interesting to search for a dynamical mechanism to allow for the diversity in a system with interacting population dynamics. Ikegami and the author have studied a population dynamics model with interaction among species, mutation, and mutation of mutation rates [23, 24]. In particular, a model with interaction among hosts and parasites has been studied. Each species is coded by a bit sequence as in §4, whose fitness has a rugged or flat (neutral) landscape. The interaction between a host and a parasite is assumed to depend on the Hamming distance between their bit sequences. When the interaction between hosts and parasites is weak, the mutation rates of species decrease with evolution. The dynamics of the whole species is reduced to a direct product of isolated sets of host-parasite population dynamics. When the interaction is strong, on the other hand, mutation rates are sustained at a high level, where many species form a network of population dynamics. This network consists of species connected by single point mutations. Many species are percolated in the gene space. Note that this network is dynamically sustained. The population of each species oscillates chaotically in time. The oscillation is high-dimensional chaos with small positive Lyapunov exponents. (“High-dimensional” here means that the number of positive exponents is large). If the mutation rate were zero, the dynamics of each species would be essentially disconnected. Then some host-parasite pairs would show strong chaos, while others would show periodic or fixed point dynamics. By sustaining a high mutation rate, chaotic instability is shared by almost all species, leading to weak high-dimensional chaos. By the term “weak”, we mean that the maximum Lyapunov exponent is close to zero, and that the amplitude of oscillation of each species is small. Our system has a tendency to evolve towards such weak, high-dimensional chaos. Here we propose a conjecture that diversity in an evolutionary system with interaction of many replicating units maintains its dynamical stability by forming a weak high-dimensional chaotic state, rather than in a fixed point or in strong chaos. We have coined the term homeochaos for this homeo-
dynamic state. The following three points capture the essence of homeochaos: (i) **Weak Chaos**: Homeochaos suppresses strong chaos. The maximal Lyapunov exponent is positive, but is close to zero. The oscillation amplitude is not large. This weak chaos, for example, is essential to avoid an overly violent change or extinction in the population dynamics. (ii) **High-Dimensional Chaos**: Homeochaos is high-dimensional chaos. There are many positive Lyapunov exponents, although their magnitude is small, and there are many degrees of freedom. (iii) **Dynamic Stability and Robustness against External Perturbations**: Homeochaos provides dynamic stability for a complex network. The robustness of homeochaos is easily seen by introducing an external perturbation to the population dynamics. If the population dynamics follow low-dimensional chaos, the amplitude of population change is sometimes very large. The oscillation can bring about a state with very small population size (see Fig.4). When an external perturbation is applied at this time, the number of population may go to zero. On the other hand, the oscillation amplitude is small in homeochaos (see Fig.4), since the chaos is very weak. Thus populations of species fluctuate around some value far from zero. Hence species are not easily driven to extinction by external perturbations. The above three features are strongly interrelated. The stability and robustness (iii) are sustained by the suppression of strong instability given by (i). By (ii), strong chaotic instability is shared by many modes, implying the weak chaos per degrees of freedom (the point (i)). The point (i) is a feature common with homeochaos and the edge of chaos. However homeochaos is not sustained at a critical point, but is more robust against a parameter change. Also the degrees of freedom are not discussed in the edge of chaos, but they are essential to homeochaos. Remnants of clustering of oscillation are important in sustaining homeochaos. Indeed the chaotic itinerancy seen in clustering is sometimes seen in homeochaos. The oscillation of some populations of some species form partial clustering over some time steps. The connection between homeochaos and clustering may not be so surprising. If chaos were too strong, oscillations of many elements would not keep any relationship, and they would become completely desynchronized. If chaos were completely suppressed, clustering with few number of clusters would often follow. To keep weak and high dimensional chaos, partial clustering with chaotic itinerancy is the most preferable state. Homeochaos in the formation of networks will be important in various levels of biological networks. Maybe the most straightforward application will be found in immune networks and quasispecies of viruses. In the immune system, antibody-antigen interactions are similar to host-parasite interactions. An antigen is damaged by ”matched” antibodies. An antibody itself is damaged by a different class of antibodies, as Jerne [27] proposed. High mutation rates are sustained for antibodies, and the concentrations of antibody species oscillate in time. We note that these features are shared in common with our model above and its numerical results. Correspondingly, viruses keep high mutation rates, and the population dynamics maintains high mutation rates. The population dynamics of quasispecies [26] is of interest from the viewpoint of homeochaos. Indeed the term “quasispecies” of [26] corresponds to our term “meta-species”, a symbiotic network connected by mutation in [23]. At a more macroscopic level, metabolic oscillations of cells may form homeochaotic dynamics, as will be discussed in the next section. At a further macroscopic level, physiology and medicine are the birthplace of the term “homeostasis”. However, data have recently been accumulating showing that the healthy state is not one of ”stasis”, but rather exhibits an irregular temporal dynamics. So far, it is hard to conclude that the dynamics of the healthy state is chaotic. Possibly this difficulty may be due to the high-dimensionality of the dynamics, where no powerful means of diagnosis from the data is available. Indeed, low-dimensional chaos has to date been found to be associated with unhealthy states of
the heart rhythm and EEG. One possible conjecture is that a healthy state is sustained by homeochaos rather than homeostasis, since homeochaotic dynamics is neither too irregular nor too regular. At the most macroscopic level, an ecological network can be a candidate for homeochaos. As Elton has discovered in the forest of England [28], an ecological system with diversity is robust against external perturbations. A typical example of a complex ecological network is found among the species in a tropical rain forest. The ecology there consists of a huge number of species, whose population size is often very small. So far, the dynamics of the population of species in a complex ecological system has not been seriously studied, but the ecology is believed to be in a dynamic state, not at a stationary state. This diversity and dynamics are also seen in our population dynamics showing homeochaos. It is strongly hoped that the population dynamics of rain forest species is measured soon. When this is done, we believe the dynamics will be found to be homeochaotic. We also believe that the mutation rate itself will be found to be larger than the normal level. If this is the case, we may assume that the coupling among species is effectively larger than in temperate zones. We may also hope that our homeochaos is important in a complex network system in general; for a dynamical network system with many units evolving according to some inherent dynamics. Such examples may include neural systems, computer networks, economics, and sociology. Our homeochaos provides a key principle for the formation of cooperation and the dynamic stability required in such systems.

6 Source of Novelty and Growth of Diversity: Open Chaos

In the previous section we have studied the maintenance of diversity. How about its creation? In §3, we have discussed a possible theoretical basis for the origin of diversity. The most typical origin of diversity is seen in cell differentiation. The formation and maintenance of a society of differentiated cells [29] is also important for the origin of multicellular organism. By cell division, each cell reproduces itself with differentiation and forms a network of cell society. Is the dynamic clustering mechanism in §3 relevant to cell differentiation? To consider this problem, quite remarkable experimental results are reported [30]: E-coli with identical genes can split into several groups with different enzymatic activities. Even prokaryotic cells with identical genes can be differentiated there. Furthermore, cells are under liquid culture, thus they are in an identical environment. This experimental result must be surprising to molecular biologists and also to those who study differentiation in the context of spatial pattern formation (e.g., along the line of Turing instability): Neither genetic nor spatial information is essential to the differentiation in the experiment. On the other hand, the experimental result may not be so surprising in the light of the dynamic clustering discussed in §3. Nonlinear metabolic reaction is involved in each cell, as well as nonlinear interaction with the soup. Cells interact globally with all other cells through the soup. Thus it is possible to expect that cell’s chemical oscillations differentiate and form some groups. Recently Yomo and the author proposed a model of cell differentiation based on the idea of dynamic clustering in §3 [31]. The model consists of metabolic reaction, active transport of chemicals from a medium (soup), and cell division. The metabolic reaction is nonlinear due to the feedback mechanism of catalytic reaction. The interaction of cells is global, due to the competition for taking chemicals (resources to produce enzymes and DNA) from the medium. Cell division is assumed to follow the chemical activities in each cell. The division speed of
a cell is assumed to be proportional to its average chemicals included therein. From a dynamical systems point of view, the model has a novel feature, not included in the globally coupled map. Here the number of degrees of freedom varies via cell division. When a new cell is born, we need additional degrees of freedom to indicate the cell’s state. Numerical results of the model show that there are three successive stages in the growth of the number of cells: coherent growth, dynamic clustering, and fixed cell differentiation (see Fig. 5). At the first stage, oscillations of chemicals of all cells are synchronized, and they divide synchronously. Hence the number of cells increases as 1, 2, 4, 8, 16, · · ·. The second stage starts when the oscillations of chemicals lose their synchronicity, and the dynamic clustering of §3 is observed. At the last stage, some (active) cells start to have more chemicals than others. Disparity in chemical activities is observed. The speed of division of active cells is at least $10^2$ times faster than other cells. The active cells may correspond to germ cells, while others correspond to somatic cells. Here, somatic cells are also differentiated according to the concentration of contained chemicals. The oscillations of chemicals at the third stage sustains weak chaos. We note that the third stage is robust against external perturbations. Since all cells compete for resources from the medium, the above results may be interpreted as follows. At the second stage, cells form a time sharing system for resources, by the clustering of oscillations, while the differentiation between poor and rich cells is formed spontaneously at the third stage. The emergence of the third stage is rather new, unexpected from the dynamic clustering in §3. Indeed, we may introduce a new concept, which we call “open chaos.” We propose open chaos as a novel and general scenario for systems with growing numbers of elements. By the active transport dynamics of chemicals, the difference between two cells can be amplified, since a cell with more chemicals is assumed to get even more. Tiny differences between cells can grow exponentially if parameters satisfy a suitable condition. A grown cell is divided into two, with an (almost) equal partition of the contained chemicals. This process looks quite similar to chaos in the Baker’s transformation, involving stretching (exponential growth) and folding (division). One difference between our cell division mechanism and chaos is that phase space itself changes after a division in the former, while the orbit comes back to the original phase space in the stretching-folding mechanism of chaos. Our “open chaos” concept is a novel and general mechanism of instability and irregular dynamics in a system with growing phase space. In studies of artificial life, the term “open ended evolution” often refers to a dynamics whose phase space attains more dimensions with the appearance of new species, strategies, and so forth. Open chaos provides a mechanism for the way in which chaotic instability in dynamical systems can trigger the expansion of the dimension of phase space. It is interesting to extend the present open chaos to areas studied in connection with open ended evolution, such as economics, sociology, game theory, and so on.

7 Beyond Top-down and Bottom-up Approaches

There have been long debates between top-down and bottom-up approaches in artificial intelligence and neural networks. In the bottom-up approach some kind of “order parameters” constructed from a lower level gives a higher level, related with some macroscopic behavior. In the top-down approach, a few instructions send messages to lower-level elements. Of course it is possible to include a weak feedback between top/bottom levels, starting from each approach. An example is a simulation of ants with pheromone. The dynamics of lower-level units (ants’ motion) leads to a collective field of pheromone, which
governs the motion of the lower-level units. Since the dynamics of lower-level units is
governed by the higher-level dynamics, this scheme is essentially analogous with the Pri-
gogine’s dissipative structure [32] or Haken’s slaving principle [33]. In these approaches, it
is assumed that the top level is represented by small number of degrees of freedom, while
the bottom level may involve a huge number of degrees of freedom. Furthermore, relation-
ships between elements are fixed. Although it is possible to include a nontrivial dynamics
at a macroscopic level, the behavior of each element is passive and totally susceptible to
a higher-level. Our network of chaotic elements provides a different mechanism, in the
following sense: first, the top level is not necessarily represented by only a few degrees
of freedom; second, relationships between elements at the lower level can dynamically
change; third, elements are not passive, but are active and dynamical [34]. The first
point may seem just a complication at first glance, but this is not necessarily so. Often
motion governed by few degrees of freedom emerges, which however, does not last forever
due to the second point (dynamic change of relationships). Again high-dimensional mo-
tion comes back, until another structure emerges. This mechanism, described as chaotic
itinerancy in §3, is essential to replace the top-down and bottom-up approaches. In a net-
work of chaotic elements, for example, the order at the top level is destroyed by chaotic
revolt against the slaving principle [3], in contrast with passive elements in traditional
approaches. In the population dynamics model by Ikegami and the author (mentioned in
§5), the higher-level corresponds to the survival of species as a collection of species. The
higher level emerges from the bottom level, but it is not necessarily represented by few
degrees of freedom.

8 Conclusion

Summing up, we have discussed a chaotic scenario of evolution, which allows for the
dynamical change of relationships of units, and spontaneous formation and destruction
of upper levels. The origin and maintenance of complexity and diversity are explained
through this scenario. We hope that chaos can remain a source of complexity, novelty,
and diversity in the studies of artificial life.

9 Acknowledgements

I am grateful to Ichiro Tsuda, Takashi Ikegami, Tetsuya Yomo, Yukito Iba, Junji Suzuki,
Kei Tokita, and Tomoyuki Yamamoto for continual criticism and stimulating discussions.
I would also like to thank Chris Langton for critical reading of the manuscript, and
discussions on his collectionism. The work is partially supported by Grant-in-Aids for
Scientific Research from the Ministry of Education, Science, and Culture of Japan.

References

[1] D. R. Hofstadter, Gödel, Escher, Bach, 1979, Basic Books

[2] M. Kimura, The neutral theory of molecular evolution, Cambridge Univ. Press, 1983

[3] N.H. Packard, in Dynamic Patterns in Complex Systems, eds. J. Kelso, A.J. Mand-
dell, M.F. Shlesinger, World Scientific, p. 293-301 (1988). See also the paper by M.
Mitchell et al., Physica D (1994), to appear (Special issue of Oji Seminar Proceeding), however.

[4] P. Grassberger, Int. J. Theor. Phys. 25 (1986) 907
[5] J.P. Crutchfield and K. Young, Phys. Rev. Lett., 63 (1989) 105
[6] C. Langton, Physica 42D (1990) 12
[7] S.A. Kaufmann and S. Johnson, J. Theor. Biol. 149 (1991) 467
[8] K.Kaneko, Physica 34 D (1989) 1
[9] J. Holland, Escaping Britleness in Machine Learning II, eds. R.S. Mishalski, J.G. Carbonell, and T.M. Miechell (Kaufman, 1986)
[10] K. Kaneko and J. Suzuki, “Evolution to the Edge of Chaos in Imitation Game”, Artificial Life III (1993), in press; “Imitation Games”, to appear in Physica D (1994), as a special issue of Oji seminar.
[11] C.K. Catchpole, Vocal communication in birds, Edward Arnold Press, 1979
[12] Krebs J.R., Ashcroft,R., and Weber,M, Nature 271 (1978) 539-42
[13] R. May, Nature 26 (1976) 459; P. Collet and J.P. Eckmann, Iterated Maps on the Intervals as Dynamical Systems, Birkhasuer, Boston-Basel-Stuttgart, 1980.
[14] J. Maynard Smith, Evolutionary Genetics, Oxford Univ. Press, 1989
[15] K. Kaneko, Phys. Rev. Lett. 63 (1989) 219; Physica 41 D (1990) 137; Physica 54 D (1991) 5; J. Phys. A, 24 (1991) 2107
[16] K. Kaneko, Phys. Rev. Lett. 65 (1990) 1391; Physica 55D (1992) 368
[17] It is equivalent to $y_{n+1}(i) = f[(1-\epsilon)y_n(i) + \frac{\epsilon}{N} \sum_{j=1}^{N} y_n(j)]$, with the aids of transformation $y_n(i) = f(x_n(i)).$ In this form, one can see clear correspondence with neural nets: If one chooses a sigmoid function (e.g., $tanh(\beta x)$) as $f(x)$ and a random or coded coupling $\epsilon_{i,j}$, a typical neural nets is obtained.
[18] I. Tsuda, “Chaotic Neural Networks and thesaurus”, in Neurocomputers and Attention, (eds. A.V. Holden and V. I. Kryukov, Manchester Univ. Press, 1990)
[19] K.Ikeda, K. Matsumoto, and K. Ohtsuka, Prog. Theor. Phys. Suppl. 99 (1989) 295
[20] Kruger, J., eds. 1991, Neuronal Cooperativity, New York, Springer-Verlag
[21] K. Kaneko, “Hypercubic Coupled Maps”, to be submitted to Physica D (1992)
[22] Here we use an identical dynamics for all species. To include the fitness, one can adopt element-dependent dynamics $f_i(x)$, instead. In this case we can see some effects of clusterings, too.
[23] K. Kaneko and T. Ikegami, “Homeochaos: Dynamics Stability of a symbiotic network with population dynamics and evolving mutation rates” Physica 56 D (1992), 406-429
[24] T. Ikegami and K. Kaneko, “Evolution of Host-Parasitoid Network through Home-chaotic Dynamics” Chaos 2 (1992) 397-408

[25] R. May, Stability and Complexity in Model Ecosystems, Princeton Univ. Press (1973)

[26] M. Eigen, J. McCaskill, and P. Schuster, Journal of Physical Chemistry, 92 (1988) 6881; M. Eigen, Sci. Am. (1993) July, 32

[27] N.K. Jerne, in The neurosciences: A study program, eds. G.C. Quarton, T. Melnechuk, and F.O. Schmitt, Rockefeller Univ. Press

[28] C.S. Elton, The Pattern of Animal Communities, 1966, Methuen and Co.Ltd. London

[29] L.W. Buss, Evolution of Individuality, Princeton Univ Press, 1988

[30] E. Ko, T.Yomo, I. Urabe, to appear in Physica D (1994), as a special issue of Oji seminar.

[31] K. Kaneko and T. Yomo, “Cell Division, Differentiation, and Dynamic Clustering”, to appear in Physica D (1994), as a special issue of Oji seminar.

[32] G. Nicolis and I. Prigogine, “Self-organization in Nonequilibrium Systems”,

[33] H. Haken, “Synergetics”, Springer, 1978

[34] The necessity of active elements at low levels and synaptic changes of higher-levels are also emphasized as “collectionism”, by C. Langton (private communication). Our approach shares these features with collectionism. We have explicitly shown that chaos can provide a mechanism for these features, here.
Figure Caption Fig.1: Emergent landscape: Average score for the players with parameters within \([a_i, a_i + \Delta]\) is plotted with \(a_i = -1 + i \times \Delta\). \(t_{trn} = 255\). The number of players is fixed at 200. a) The mutation rate \(\mu = 0.1\), and \(T = 32\). The bin size \(\Delta = 0.001\). Sampled for time steps from 1000 to 1500, over all players. (See for details \[1\], from which the figure is adapted). b) \(\mu = 0.001\), and \(T = 128\). The bin size \(\Delta = 0.002\). Sampled for time steps from 750 to 1000, over all players. Fig.2: Average score of the game vs. Lyapunov exponents. Simulation is carried out with \(\mu = 0.05\), \(t_{trn} = 255\), and \(T = 128\) by fixing the population of birds at 200. Average scores are obtained from the histogram of Lyapunov exponents, for which we use a bin size of 0.01 for \(-1 < \lambda < 1\), while it is set at 0.1 for \(\lambda < -1\) (since the sample there is rather sparse). Sampled over time steps from 500 to 750 over all players (whose number is fixed at 200). Fig.3: Space-time diagram for the coupled map lattice on a hypercubic lattice with \(k = 5\) (i.e., \(N = 2^5\)). For local dynamics \(f(x) = 1 - 1.52x^2\) is adopted, while the coupling strength \(\epsilon\) is set at 0.3. On the corresponding pixel at a given time and element, a bar with a length proportional to \((x_n(i,j) - 0.1)\) is painted if \(x_n(i,j) > 0.1\). Every 4th time step is plotted from 10000 to 12000. Elements are aligned according to its binary representation; i.e., \(0 = 00000\), \(1 = 00001\), \(2 = 00010\), \(\cdots\), \(63 = 11111\). At the stage A, elements split into two clusters \(0 \ast \ast \ast \ast \ast\) and \(1 \ast \ast \ast \ast \ast\), while they split into \(0 * * * * *\) and \(1 * * * * *\) at the stage B, \(\ast * * * * *\) and \(\ast 1 * * * *\) at the stage C, \(0 * * * * * * * *\) and \(1 * * * * * * * *\) at the stage D, \(\ast * * * * 0\) and \(\ast * * * * 1\) at the stage E, and again into \(0 * * * * * * * *\) and \(1 * * * * * * * *\) at the stage F. Fig.4: Oscillation of total populations for hosts (solid line) and parasites (dashed line), in the model in \[23\]. Initial chaotic oscillations with large amplitudes are suppressed simultaneously with the increase of averaged mutation rate around time = 2300. Up to the time, the population dynamics is a direct product of low-dimensional chaos, and includes a violent change of populations. After the time, the oscillation amplitude is much weaker (although chaotic), where homeochaos is attained with the increase of the mutation rates. See for details \[23\]. Fig.5: Overlaid time series of a chemical at each cell. As the number of cells grows, the oscillation starts and clustering emerges at the stage II. At the stage III, disparity of the chemical is clearly seen. See for details \[31\], from which the figure is adapted.