Capturing Intracellular pH Dynamics by Coupling Its Molecular Mechanisms within a Fully Tractable Mathematical Model

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Abstract

We describe the construction of a fully tractable mathematical model for intracellular pH. This work is based on coupling the kinetic equations depicting the molecular mechanisms for pumps, transporters and chemical reactions, which determine this parameter in eukaryotic cells. Thus, our system also calculates the membrane potential and the cytosolic ionic composition. Such a model required the development of a novel algebraic method that couples differential equations for slow relaxation processes to steady-state equations for fast chemical reactions. Compared to classical heuristic approaches based on fitted curves and *ad hoc* constants, this yields significant improvements. This model is mathematically self-consistent and allows for the first time to establish analytical solutions for steady-state pH and a reduced differential equation for pH regulation. Because of its modular structure, it can integrate any additional mechanism that will directly or indirectly affect pH. In addition, it provides mathematical clarifications for widely observed biological phenomena such as overshooting in regulatory loops. Finally, instead of including a limited set of experimental results to fit our model, we show examples of numerical calculations that are extremely consistent with the wide body of intracellular pH experimental measurements gathered by different groups in many different cellular systems.

Introduction

Distribution of charges within biological molecules is crucial, not only for reactivity and catalysis, but also as it determines their solubility, their particular folding, and dictates the spatio-temporal sequence of their interactions. In this context, the pH of the solution bathing these biological molecules is a key parameter, since its value determines the protonation of the acid-base groups that are especially abundant in macromolecular assemblies. Furthermore, as many enzymes and cellular regulators exhibit a strong pH dependency, the modification of the protonation of key residues can deeply impact their function. For these reasons, genomes necessarily contain pH-dependency information, which is expressed in the proteome [1]. The complete information for intracellular pH determination is a convoluted interplay between the abundance and the distribution of protonable groups in biological molecules, their pKa values and the expression, stability, kinetic and affinity parameters of the pH regulating systems. Accordingly, providing a fully tractable model for intracellular pH regulation is a challenging problem, and several studies have been aimed at building essentially heuristic models [2–5] for intracellular pH regulation.

The past decades have witnessed the detailed molecular characterization of the protagonists that regulate the concentrations of cellular acid-base equivalents, in term of both their kinetics and the affinities for their substrates [6,7]. Significant efforts have also been invested to describe intracellular buffering mechanisms and proton diffusion in cells adequately [8,9].

Based on this, we develop here a different, bottom-up approach at the interface between biology, physics, chemistry and mathematics. We construct a model that encompasses the individual molecular mechanisms for these regulators defined by their own kinetics and by their experimentally measured microscopic parameters. This requires the inclusion of the chemical reactions between the involved reactive species. This non-empirical process guarantees the construction of a physically coherent, fully integrated and tractable model (i) for cellular proton dynamics and (ii) for steady-state pH regulation.

In the present study, we choose to keep the system simple and modular by assuming that the cell surface and volume are fixed to their average values and by using the ubiquitous Na+/H+ exchanger NHE-1 and Cl⁻/HCO₃⁻ exchanger AE2 as the main transmembrane acid-base transporters. We also include the electrical gradient generated by the Na+/K-ATPase across the membrane and the permeabilities associated to Na⁺, K⁺ and Cl⁻ background currents measured in non-excitable cells. Therefore, our model computes the distribution of the other cationic and anionic species and their variations as a function of proton concentration.
These pumps and transporters show a very high sequence conservation within different mammalian species and possess very similar constants for their substrates. Based on this, we built our model using widely accepted values from the literature even if they had been measured from different mammalian species. We will further see that this is validated by our results, which show that pH regulation is very resilient against variations of those thermodynamic constants.

It is demonstrated that our model gives (i) a robust, experiment based prediction of the temporal evolution of the pH, (ii) a simple analytical value for its steady state, (iii) all the other ionic concentrations related to the proton regulation, (iv) and a reduced differential equation for describing the full pH balance.

This enables the testing of biologically-relevant situations whilst discriminating between critical parameters and rate limiting steps toward those factors that can be widely changed with virtually no effect on cellular homeostasis.

Methods

Datasets used for the Model

We report most of thermodynamical data, the common ionic environments, and the justification of the kinetic equations in the Datasets in File S1. In the following, we illustrate the specific behavior of the involved physical, chemical or biological components.

Ionic Flows and Potential through the Membrane

Let us depict the cellular model represented in Figure 1 mathematically. We assume that the cell geometry is fixed by neglecting that water flows through the membrane. The charge balance is controlled by passive, electroneutral, electrogenic flows and capacitive currents that are described as follows.

Passive flows. If X represents a chemical species in Figure 1, with an inner concentration [X] and an outer concentration [X]out, then it flows out of the cell through the membrane surface S due to a permeability P X(ζm). Here, ζm = FEm/(RT) represents the reduced electric potential where F is the Faraday constant, R is the molar gas constant, T is the absolute temperature, and Em is the electric potential difference between the cytosol and the outer medium. The Goldman-Hodgkin-Katz flux equation [10] provides the outward molar flux jX as

\[ j_X = -P_X(\zeta_m)\Psi(\zeta_m, \zeta_{m})[X]_{out} - [X]_{in}e^{\zeta_m} \]

with \( \Psi(u) = u(e^u - 1)^{-1} \) and \( \zeta_X \) is the algebraic charge of X. The associated passive outward electric flux is \( j_X = \zeta_X F X \) and the whole cell passive outward electric current is \( I_X = S j_X \). We can simply convert the flux into an intake molar rate for a given cell volume \( V \) as

\[ \hat{\varepsilon}_X[X]_{\text{passive}} = -\frac{S}{V} j_X = \zeta_X(\zeta_m - \zeta_{X})[X]_{out}/[X]. \]

For the cellular system used in the electrophysiological measurements we recorded significant currents only for \( K^+ \), \( Na^+ \) and \( Cl^- \) (CCL39 cells, see Figure S1 in File S1). This allows the determination of the corresponding permeabilities. Any other species can be taken into account if other cells are considered, and if values are available or measurable.

Electroneutral transporters. The electroneutral AE2 exchanger keeps \( Cl^- \) ion concentration above its Nernst potential [11], and thus is assumed to work in the forward direction. AE2 is then operating with a Hill mechanism [12], inducing a whole cell exchange rate

\[ \hat{\varepsilon}_X[Cl^-]_{AE} = -\hat{\varepsilon}_X[\text{HCO}_3^-]_{AE} \equiv \rho_{AE} \]

with \( \rho_{AE} = V_{AE} [\text{HCO}_3^-]/(K_{AE} + [\text{HCO}_3^-]) \) and where \( V_{AE} \) is the cellular maximal \( \text{HCO}_3^-/Cl^- \) exchange rate, and where \( K_{AE} \) is the bicarbonate affinity. Unless indicated, we will use a Michaelis-Menten behavior, namely \( \gamma = 1 \), and \( K_{AE} \) is about 10 mM.

We use the established mechanism [13] of the NHE-1 exchanger that results in the whole cell exchange rate

\[ \hat{\varepsilon}_X[Na^+]_{NHE} = -\hat{\varepsilon}_X[\text{H}^+]_{NHE} \equiv \rho_{NHE} \]

with \( \rho_{NHE} = V_{NHE}\sigma_{NHE}([\text{H}^+]_t)/K_t \) and where \( V_{NHE} \) is the cellular maximum \( \text{H}^+/Na^+ \) exchange rate, \( K_t \) is about 1.8·10^{-8} M and

\[ \sigma_{NHE}(x) = \frac{x(1+x)+L_0 C_x(1+C_x)}{L_0^2(1+C_x)^2+(1+x)^2}. \]

with \( C = K_t/K_n \), \( K_t \approx 3.6·10^{-6} \) M and \( L_0 \approx 10^3 \).

Any other electroneutral transporter could be similarly described and therefore inserted into the model.

Electrogenic currents. We restrict ourselves to the sodium-potassium pump that exchanges three inner \( Na^+ \) with two outer \( K^+ \) according to

\[ \frac{1}{2} \hat{\varepsilon}_X[\text{K}^+]_{Nak} = -\frac{1}{3} \hat{\varepsilon}_X[\text{Na}^+]_{Nak} \equiv \rho_{Nak}(\zeta_m). \]

where \( \rho_{Nak}(\zeta_m) = V_{Nak}\sigma_{Nak}(\zeta_m, [Na^+]_t) \). \( V_{Nak} \) is the cellular maximum \( Na^+/K^+ \) exchange rate, and we combined the experimental data found in published studies [14,15] to estimate (see Datasets in File S1)

\[ \sigma_{Nak}(\zeta_m, [Na^+]) = \frac{[1 + \tanh(0.39\zeta_m + 1.28)] [Na^+]_t}{K_{Nak} + [Na^+]_t}. \]

with \( K_{Nak} \approx 10^{-15} \) mM.

Electric potential evolution. The cytosol and the outer medium must remain globally electroneutral. Conversely, charge accumulation polarizes the membrane due to its surface capacitance \( C_m \approx 1 \mu F/cm^2 \) (see Materials and Methods in File S1). The total capacitance of cell membrane is \( C = SC_m \).

We take into account both the passive and electrogenic actors, respectively defined by equations (1) and (5), which are involved in the electric potential regulation. This results in

\[ C \hat{\varepsilon}_m E_m + \sum_X I_X(E_m) + I_{Nak}(E_m) = 0, \]

with the electric conversion \( I_{Nak} = F V \rho_{Nak} \).

Chemical Physiology

So far in our modeling, the species passing through the membrane are \( \text{H}^+, \text{HCO}_3^-, \text{Cl}^-, \text{K}^+ \) and \( Na^+ \). Obviously, the
first two of them are directly involved in the set of protic reactions that govern the pH. Since the physiological range of pH lies around 7, we must monitor in our analysis the self-ionization of water

$$H_2O \rightleftharpoons H^+ + HO^-, \quad K_w = [H^+][HO^-]. \quad (7)$$

Then, we include the three components of the carbonated system. The partial pressure $\Pi_{CO_2}$ of carbon dioxide equilibrates with aqueous $CO_2$ (according to the Henry law), which undergoes two consecutive dissociations. Those reactions and their equilibrium constants are summarized below:

$$CO_{2gas} \rightleftharpoons CO_{2aq}, \quad K_H = \Pi_{CO_2}/[CO_2] \quad (8)$$

$$CO_{2aq} \rightleftharpoons H^+ + HCO_3^-, \quad K_1 = [H^+][HCO_3^-]/[CO_2] \quad (9)$$

$$HCO_3^- \rightleftharpoons H^+ + CO_3^{2-}, \quad K_2 = [H^+][CO_3^{2-}]/[HCO_3^-]. \quad (10)$$

Since the dissolved $CO_2$ is very unstable in water, and especially in presence of physiological carbonic anhydrase [16], we shall directly merge the equilibria (8) and (9) to get rid of $CO_{2aq}$ and obtain an equivalent equilibrium:

$$CO_{2gas} \rightleftharpoons H^+ + HCO_3^-, \quad K'E = K_1/K_H = [H^+][HCO_3^-]/\Pi_{CO_2}. \quad (11)$$

Finally, we model the protic behavior of all the other species within the cytosol by a single equivalent buffer that we name $Y$, acting as

$$HY \rightleftharpoons H^+ + Y^-, \quad K_Y = [H^+][Y^-]/[HY] \quad (12)$$

with $Y_0 = [HY] + [Y] \approx 50-70 \text{ mM}$ (see [8] for details). Additionally, we do not consider the diffusive effects within the cell, by assuming homogeneous ionic concentrations. Similarly, the outer...
media is considered as an infinite bulk with constant physiological concentrations of the different entities under consideration.

**Results**

**General Theoretical Results**

**Self-consistency and modularity.** To be self-consistent, our model must ensure that each ion and the membrane potential are sufficiently maintained by physico-chemical processes (enzymatic transport or chemical reaction) in order to avoid some non-physical accumulation or discrepancy such as negative concentration. Accordingly, we propose a model for a generic cell, restricted to the previous components (transporters, ionic permeabilities and chemical reactions) which fulfills these criteria. Noticeably, any other effector that acts redundantly for pH regulation or is expressed in specialized cells can be implemented, provided that the above self-consistency is preserved. As an example, we will show further how to handle the lactate/\(H^{+}\) production and transport. The same methodology applies to other mechanisms (such as \(H^{+}\)-ATPase or \(Na^{+}\)-coupled-bicarbonate transporters) and any additional chemical reaction.

**Full formal system dynamics.** For any reaction indexed by \(k\), we will note \(\gamma_{k}\) its associated rate, which is the derivative of its chemical extent with respect to time. For each chemical species, the concentration temporal derivative is the appropriate summa-

\[
\begin{align*}
\mathcal{C} \dot{E}_{m} &= \mathcal{F}(\lambda_{K^{+}} + \lambda_{Na^{+}} - \lambda_{Cl^{-}} - \rho_{NaK}) \\
\dot{\gamma}_{i}[K^{+}] &= \lambda_{K^{+}} + 2\rho_{NaK} \\
\dot{\gamma}_{i}[Na^{+}] &= \lambda_{Na^{+}} - 3\rho_{NaK} + \rho_{NHE} \\
\dot{\gamma}_{i}[Cl^{-}] &= \lambda_{Cl^{-}} + \rho_{AE} \\
\dot{\gamma}_{i}[HO^{+}] &= \gamma_{H^{+}} \\
\dot{\gamma}_{i}[CO_{3}^{2-}] &= \gamma_{H^{+}} \\
\dot{\gamma}_{i}[H^{+}] &= -\rho_{AE} + \lambda_{H^{+}} - \gamma_{H^{+}} \\
\dot{\gamma}_{i}[HCO_{3}^{-}] &= -\rho_{NHE} + \lambda_{H^{+}} + \lambda_{H^{+}} + \gamma_{H^{+}} \\
\end{align*}
\]

where \(\gamma_{i}, \lambda_{i}, \gamma_{Y}, \gamma_{X}\) are respectively the molar rate of the water ionization (7), the direct formation of \(HCO_{3}^{-}\) (11), the dissociation of \(HCO_{3}^{-}\) (10), and the deprotonation of the equivalent buffer (12). The first equation of the above system is rewritten from relation (6).

The main characteristic of protic reactions in water is that they have very short relaxation times, from a few microseconds for \(CO_{2}\) [17] to a few milliseconds for \(CO_{2}\) with carbonic anhydrase [16] or without carbonic anhydrase [18]. Since the transmembrane exchanges of protic species (through \(NHE = I\) or \(AE2\)) are expected to have characteristic times much larger, we can consequently make the assumption that each proton equilibrium is in fact a fast pre-equilibrium. It follows that each involved reaction quotient always matches its corresponding thermodynamical reaction constant; this can be seen as a set of constraints applied to the chemical composition of the aqueous solution. Consequently, if we want to impose a perturbation of this composition then those pre-equilibria shall instantaneously produce the mandatory chemical extents which ensure that the final composition respects the chemical constraints. Accordingly, the thermodynamical knowledge of the equilibria (7) to (12) is sufficient to solve the kinetic equations (13) in this particular biological context.

In the following, we detail the treatment of the protic reactions rates within the equations (13). We show (see Methods in File S1) how to derive a set of reduced differential equations for the 5 dynamic variables \(E_{m}, [K^{+}], [Na^{+}], [Cl^{-}]\) and \([H^{+}]\), within this pre-equilibria approximation.

**Steady state characterization.** The steady-state values, which we note within asterix, are obtained by setting to zero the temporal evolution in the differential system (13) leading to

\[
\begin{align*}
0 &= \gamma_{K^{+}}^{*} + \gamma_{Na^{+}}^{*} - \gamma_{Cl^{-}}^{*} - \rho_{NaK}^{*} \\
0 &= \gamma_{K^{+}}^{*} + 2\rho_{NaK}^{*} \\
0 &= \gamma_{Na^{+}}^{*} - 3\rho_{NaK}^{*} + \rho_{NHE}^{*} \\
0 &= \gamma_{Cl^{-}}^{*} + \rho_{AE}^{*} \\
0 &= \rho_{NHE}^{*} - \rho_{AE}^{*}.
\end{align*}
\]

This system is under-determined since the sum of the first equation and of the last two ones minus the second and the third one is zero; this is expected since the evolution of \(E_{m}\) is the exact conservation of the global electric charge. In our model, the latter decomposes into an intrinsic charge \(Q\) of all the considered components and an excess charge \(Q^{\prime}\) of all the other "spectator" species (proteins, other ions…), leading to \(CE_{m} = \Delta Q + \Delta Q^{\prime}\), where \(\Delta\) means the difference between the cytosol and the outer medium values. It is the integrated form of equation (6). As a consequence, the initial condition of the differential system (13) gives \(\Delta Q^{\prime}\).

In order to determine steady-state values, we first obtain the electric equation \(0 = \frac{3}{2} \gamma_{K^{+}}^{*} + \gamma_{Na^{+}}^{*} - \gamma_{Cl^{-}}^{*}\) which is here equivalent to

\[
\gamma_{m}^{*} = \text{ln} \left( \frac{\frac{3}{2} P_{K^{+}}^{*} [K^{+}]_{out} + P_{Na^{+}}^{*} [Na^{+}]_{out} + P_{Cl^{-}}^{*} [Cl^{-}]_{out}}{\frac{3}{2} P_{K^{+}}^{*} [K^{+}] + P_{Na^{+}}^{*} [Na^{+}] + P_{Cl^{-}} [Cl^{-}]_{out}} \right)
\]

since we only deal with monovalent ions. The relation (15) is the Goldman-Hodgkin-Katz potential equation with voltage-dependent permeabilities and a potential explicitly regulated by \(Na/K\)-ATPase.

With \(V_{AE/NHE} = V_{AE}^{*} / V_{NHE}^{*}\) (see equations (3) and (4)) and \(L_{0} \gg 1\), the last equation of the system (14) can be expressed as a function of \(h^{*}\) and it reduces to a simple polynomial

\[
K_{AE} K_{r} P_{CO_{2}} \left( h^{*} / K_{r} \right)^{2} + \left( 1 - \frac{V_{AE/NHE}^{*}}{1 + \frac{1}{1 + C L_{0}}} \right) \left( h^{*} / K_{r} \right) - 1 V_{AE/NHE}^{*} = 0
\]

This analytical relation yields the steady-state pH as a function of \(P_{CO_{2}}\) and \(V_{AE}^{*} / V_{NHE}^{*}\), since \(h^{*}\) is the positive root of equation (16). The Figure 2 shows how the exact pH evolves when those parameters are changed and where the acceptable physiological limits stand. In particular, for an intracellular pH = 7.2 and \(P_{CO_{2}} = 40\ mmHg\) our model predicts \(V_{AE}^{*} / V_{NHE}^{*} \simeq 0.057\); this transport ratio matches well the experimental maximal rates of this transporters in different systems [19–21]. An interesting feature of
our model is the prediction of missing parameters (kinetic and/or thermodynamic) based on the knowledge of steady-state physiological values (see Results in File S1). For instance, a unique $[\text{H}^+]^?$ is computed from a given $V_{\text{AE}}, V_{\text{NHE}}$ and $\Pi_{\text{CO}_2}$. Conversely, the ratio $V_{\text{AE}}/V_{\text{NHE}}$ can be read on Figure 2 from the experimental measure of $[\text{H}^+]^?$ and $\Pi_{\text{CO}_2}$.

Asymptotic Kinetics Framework

**General philosophy.** The set of differential equations (13) defines a multiple-scale system (in both time and in concentrations), since it combines slow chemical rates with fast relaxing protic reactions. In either multiple-scale analysis [22] and normal forms in central manifolds [23], the slow dynamics are assessed around a stationary point. However, in the case used herein, the slow dynamics evolve on manifolds generated by the laws of mass action (corresponding to each protic pre-equilibrium established at its thermodynamical constant) and represent the only valid compositions of the system. To the best of our knowledge, this is the first time that a way to compute the constrained evolution of the all the involved concentrations has been exposed.

**Chemical system description.** If we assume that we have $N$ chemical reactions coupling $M$ species $X_{1,...,M}$, then the relevant reactions may be written in a generic form, employing the algebraic stoichiometric coefficients $v$:

$$\forall i \in 1...N, \sum_{j=1...M} v_{i,j}X_j = 0, \quad K_i = \prod_{j=1...M}[X_j]^{v_{i,j}}$$

(17)

where $K_i$ is the equilibrium constant of the $i^{th}$ reaction. We also assume that $K_i$ is temporally dependent, so as to reflect the possible variations of the external conditions (such as imposed changes in partial pressures).

**Fast pre-equilibria consequences.** We now assume that those $N$ reactions represent fast pre-equilibria. In other words, we suppose that the relaxation time of each reaction is infinitely small. Accordingly, we define the vector $\delta \bar{X}$ by its coordinates $\delta \bar{X}_i$ such that

$$\forall i \in 1...N, \delta \bar{X}_i = K_i \prod_{j \neq i}[X_j]^{-v_{i,j}} - \prod_{j \neq i}[X_j]^{-v_{i,j}} = 0.$$  

(18)

Thus, the only permissible evolutions must be satisfied at particular time $t$ and for any set of concentrations $\bar{X}$ through the relationships:

$$\begin{align*}
\dot{\delta} &= \delta \bar{X}(t, [\bar{X}]) \\
\delta &= \varepsilon(t) \bar{X} + \Phi \varepsilon(t) [\bar{X}] \\
\end{align*}$$

(19)

where $\Phi$ is the Jacobian matrix of $\delta$ with respect to $[\bar{X}]$.

**Response to perturbations.** If we perform a small modification $\delta$ of each reaction form the vector $\bar{X}$. We then obtain a modified perturbation

$$\delta \bar{X} = \delta \bar{X} + \xi^T \bar{X}$$

(20)

and this in turn must obey

$$\bar{X} = \delta \bar{X} + \xi^T \bar{X}.$$

(21)

The instantaneous chemical extent is readily computed by

$$\bar{X} = - (\Phi^T)^{-1} (\delta \bar{X} + \Phi \varepsilon(t) [\bar{X}]).$$

(22)

We can show that the $N \times N$ matrix $\Phi^T$ is invertible for any admissible set of concentrations. This purely algebraic property results from the convexity of the free enthalpies of reactions from which the expressions of $\delta$ and $\Phi$ are derived, but this purely mathematical demonstration is far beyond the scope of this article.

**Generic asymptotic kinetics.** Finally, noting that $\varepsilon(t) \bar{X}$ contains global information regarding the “slow”-changing variations of all the chemical species, the overall chemical evolution of the system is deduced from (22) by setting

$$\delta \bar{X} = \varepsilon(t) \bar{X} - \nu^T (\Phi^T)^{-1} (\Phi \varepsilon(t) [\bar{X}] + \varepsilon(t) \bar{X})$$

(23)

thereby illustrating how the fast chemical reactions are damping the slow variations.

**Numerical integration.** A modular C++ program (available upon request) was designed and exactly encodes the biological effectors, the membrane potential and the chemistry into a system of algebraically coupled numerical differential equations. The integration step was performed by an adaptive Dormand-Prince method with a fractional tolerance of $10^{-7}$.

Reduced Model for pH Dynamics

**Evaluation.** It is assumed that all the protic reactions are rapid pre-equilibria (see above), then we can derive the proton generation rate by using the previous mathematical formalism. This formalism provides an algebraic manner to decouple all the protic reactions from the catalytic ones explicitly. Implicitly however, all protic dynamics may be deduced from the evolution of $[\text{H}^+]^?$ which has unfortunately no simple expression.

**Simplified pH dynamics.** For a given cell in physiological conditions (for which the internal steady pH is around 7.2) we can neglect the presence of $\text{CO}_2^+$ and obtain a slightly simplified rate (see Methods in File S1):

$$\begin{align*}
\varepsilon(t)[\text{H}^+] &\approx \Theta([\text{H}^+]^?, Y_0) \left( \rho_{\text{AE}} - \rho_{\text{NHE}} \right) + \frac{K^?_{\text{H}^+}}{[\text{H}^+]^?} \varepsilon(t)\Pi_{\text{CO}_2} \\
\Theta &\equiv \frac{1}{[\text{H}^+]^2 + K^?_{\text{H}^+} + K^?_{\Pi_{\text{CO}_2}} + \beta_{L} Y_0} \\
\end{align*}$$

with a numerically derived value $\beta_{L} \approx 6.55 \text{ M}^{-1}$, and for a total buffer concentration $Y_0$. Here, $\rho_{\text{AE}}$ and $\rho_{\text{NHE}}$ depend only on the pH and on $\Pi_{\text{CO}_2}$, so that one can simulate the pH with only one differential equation.

**Role of the buffer.** The steady-state $[\text{H}^+]^?$ is readily recovered from equation (24). The factor $\Theta$ emphasizes the preponderant role of the chemical couplings pertaining to the evolution of the intracellular pH. Indeed, we always have $\Theta < [\text{H}^+]^?/(K^?_{\Pi_{\text{CO}_2}})$, so that for the steady state of a “normal” cell (pH $= 7.2,$
protons will not neutralize the remaining consequence, if the cell removes ousely pumped out of the cell by the regulating enzymes. As a dynamics to be even further reduced through the term in the denominator, up to 30% for a 60 mM buffer concentration.

Natural overshoot. Interestingly, our calculations predict that a vanishing physiological protic perturbation will systematically produce a pH overshoot around its steady state. Such phenomena are well known experimentally [24,25]. The exact mathematical demonstration of this phenomenon (see Methods in File S1), is valid for acidification or alkalinization, and can be applied to model other overshoots observed in different physiological regulations.

To explain this in a non-mathematical way, we may perform the following thought experiment. Let us assume that a weak protonated base HB enters the cell at its steady state. The excess of protons produced by the dissociation HB⇒H⁺ + B⁻ is continuously pumped out of the cell by the regulating enzymes. As a consequence, if the cell removes HB from its cytosol, then some protons will not neutralize the remaining B⁻. Accordingly, the initial pH is reached earlier than expected: the further removal of B straightforwardly creates an unexpected depletion of protons (basic environment) before returning to the initial situation. This describes an overshoot mechanism.

Steady-state pH: Role of Enzymatic Constants

We have investigated the changes in the steady-state pH resulting from covalent or non-covalent modification of the transporters through intracellular signaling cascades, and the effects of allosteric activators or inhibitors or mutations that affect the transporters parameters. Unless stated, we model these effects assuming that a modification is specific and affects only one thermodynamical enzymatic constant without changing the others, while we keep the kinetic ratio and PiCO₂ to their usual levels.

For NHE−1, it has been shown that within the Monod-Wyman-Changeux framework [26], the allosteric constant L₀ is modified by various stimuli such as growth factor stimulation or changes in membrane composition and tension [13,21]. This raises the question whether NHE−1 cooperativity for proton is intrinsically important for pH regulation itself and other cellular functions. Accordingly, Figure 3 depicts the results of our computations of the resulting pH following a modification of L₀. Interestingly, it has been shown widely that the activation of NHE−1 by the above-mentioned stimuli [27] decreases L₀ by one order of magnitude [13], and results in a pH increase of 0.2 to 0.3 units. Our model yields a pH increase of about 0.3 units for L₀/stim/L₀ = 0.1, which is in very good agreement with the actual experimental data. We are also able to hypothesize that in some cases NHE−1 might be regulated by altering its microscopic affinities for protons. To investigate this, we changed the Kᵣ and Kᵣ constants of NHE−1 and kept the Kᵣ/Kᵣ ratio constant, as these two affinities correspond to the same site in different conformations. Our model predicts that any important change in these microscopic affinities would produce a large pH shift, as shown on Figure 3.

For AE2, we investigated either the effect of a variation in the affinity KAE or in the Hill exponent (see Figure 3). We observe a less drastic change than those resulting from a NHE−1 modification. This makes sense from a physiological point of view: due to membrane potential and metabolic activity, cells constantly have to compensate for intracellular acidification rather than for alkalinization.

To summarize, we show that pH regulation is robust for two main reasons. Firstly, changes in the thermodynamic constants of

\[ \text{PiCO}_2 = 40 \text{ mmHg}, \] we estimate \( \Theta' < 5 \times 10^{-6} \). The protons dynamics (hence of all the protic species) are sharply damped by those chemical couplings. As expected, the buffer causes these dynamics to be even further reduced through the term \( 1 + \beta V \) in the \( \Theta \) denominator, up to 30% for a 60 mM buffer concentration.
pH regulation systems, that in vivo could arise from mutations or from interspecific variations, induce minute modifications of the steady-state pH. Secondly, pH can relax back to its physiological value, because changes in constants are very easily overcome by slight modifications of the maximal rates of the transporters, i.e. the amount of transporters expressed at the plasma membrane. This a posteriori validate the hypotheses and choices described in the introduction.

Characteristic Time Scales in pH Regulation

In order to investigate the dynamics, we integrated the differential system (13) as described previously. We choose to approximate the average fibroblast shape using a prolate spheroid model of length $25 \mu m$ and a diameter of $10 \mu m$, leading to a surface $S = 653 \mu m^2$ and a volume $V = 1309 \mu m^3$.

We found an anionic charge excess of $\Delta Q^{\alpha} \approx -30 \text{ mM}$ for this configuration, which mainly corresponds (i) to the excess of negative charges found on the surface of intracellular proteins and (ii) to the bulk of negative charges provided by the first dissociation of phosphate groups [8].

Relaxation times around the steady-state values. We performed the linear stability analysis of the differential system (13) which also provides the relaxation constants of the independent variables. We deduced the raw and typical relaxation time constants for our cell model by setting the equivalent buffer concentration to zero. Firstly, we obtain a 3 ms characteristic time that predominantly corresponds to the relaxation of $E_{m}$: obviously the membrane potential adjusting itself very quickly to a change in the ionic composition, but since it does not produce chemical species per se, it does not influence the chemical rates.

Secondly, we have two similar time constants representing the relaxation of a perturbation of all the ions within 8 and 15 minutes.

Finally, for a fixed $\Pi_{\text{CO}_2}$, the perturbed concentration $\delta[H^+]$ dynamics obeys to

$$\dot{\delta[H^+]} = -\delta[H^+] / \tau_h$$

with $1 / \tau_h = \left[ \Theta \left( \frac{\rho_{\text{nHE}} - \rho_{\text{AE}}}{\delta[H^+]} \right) \right]^{\nu}$. Since the relation (24) provides $\Theta^{\nu} \approx 5 \times 10^{-6}$, the proton relaxation time $\tau_h$ reaches about 5 minutes, which is consistent with the experimental observations from a plethora of reports [24,28,29].

Illustration on Pathophysiological Situations

Forced acidosis: $\Pi_{\text{CO}_2}$ and intracellular buffer. We simulated an artificial increase of $\Pi_{\text{CO}_2}$ of 20% during 1 minute (from 40 mmHg to 48 mmHg), followed by a return to the normal within 5 minutes. The resulting pH dynamics with and without the equivalent buffer $HCO_3^- / Y^-$ are presented in Figure 4A. As shown in Figure 4B, the $HCO_3^-$ excess increases the $Cl^-$ intake, and the acidification increases the $Na^+$ intake, while the $K^+$ level is remarkably stable as expected. The net ionic currents produce a concurrent tiny 2 µV depolarization. We note that, via the chemical couplings, the aprotic species dynamics are also dampened by the presence of the buffer. As expected, the different timescales are also respected once the perturbation is over and all the concentrations are relaxing towards their steady-state value: the pH needs only a few minutes to recover its physiological level, while the other ions rather require tens of minutes to reach their final balance. The numerical simulation also points out the predicted overshoot with and without buffering.

If $\Pi_{\text{CO}_2}$ is held to at its maximum increase, then the pH converges to a new value that can be deduced from Figure 2. In such a case, the pH curve is similar to Figure 4A except that it converges monotonously towards almost 7.18 after the initial decrease, and no overshoot occurs. At the same time, the other
ions find a different balance. The results of this specific $\Pi_{CO_2}$ jump are shown on Figure S2A&B in File S1.

**High-flow lactic ischemia.** Here we show how to expand the model in order to probe the consequences of a slight hypoxia without ATP depletion. We model it with a lactic acid production, while we hold the enzymatic constants and $\Pi_{CO_2}$ to their normal values. Since our model is modular, we first consider the dissociation of the lactic acid $\text{LaH}$, namely $\text{LaH} \rightleftharpoons \text{La}^- + \text{H}^+$ with $K_a \approx 10^{-3.86}$. The cell removes lactates and their accompanying protons (1:1) through monocarboxylate transporters (MCT) [30,31]. The latters follow a Michaelian law defined by $K_m \approx 30$ mM and by an observed maximum rate $x_a, \approx 1$ mM/min [32]. The global lactic acid production is around $x_a, \approx 1$ mM/min in an hypoxic skeletal muscle [33]. Consequently, we simply have (i) to append

\[
\begin{align*}
\frac{d}{dt}[\text{LaH}] & = -x_a + \rho_a \\
\frac{d}{dt}[\text{La}^-] & = +x_a
\end{align*}
\]

where $x_a$ is the molar rate of the lactic acid dissociation and (ii) to include $x_a$ to $\hat{e}_i[\text{H}^+]$ within the differential system (15). Here we impose

\[
\begin{align*}
\rho_a & = x_+ \\
\rho_a & = -x_- [\text{La}^-]/(K_m + [\text{La}^-]) \quad \text{if } t < T_a \\
\rho_a & = 0 \quad \text{if } t \geq T_a
\end{align*}
\]

where $T_a$ is the ischemia duration. The resulting pH is shown in Figure 5A.

Consequently, an extra term appears in the overall protonic rate:
The acidification of the cytosol produces a massive Na⁺ overload (see Figure 5B), which is experimentally observed [28] and corresponds to a stimulated action of NHE−1. For a fixed \( \Pi_{\text{CO}_2} \), the fall of \( [\text{HCO}_3^-] \) during the lactate production decreases the chloride intake, as shown in Figure 5B. The net ionic currents induce a small 200 μV hyperpolarization over the simulation. As demonstrated in the Methods of File S1, the structure of equation (28) leads to an expected pH overshoot, which occurs after \( t > T_{\text{sh}} = 15 \) minutes. The fast regulating couple formed by NHE−1 and AE2 allows the pH to follow the rate limiting lactate expulsion.

\[
\hat{e}_i[H^+] \propto (K_a + [H^+])(\rho_{\text{NHE}} - \rho_{\text{AE}}) + K_d \rho_d. \tag{28}
\]

**Discussion**

Adequacy with Experimental Data

As previously stated, the main purpose of this study was to build a mathematical depiction of intracellular pH regulation and investigate whether it had analytical solutions and produced biologically relevant simulations. This last section intends to further challenge our study by confronting real experimental data. To avoid potential biases, we decided against the use of our own data and instead to choose one of the pioneer experiments within the large body of published intracellular pH measurements generated by independent groups in the last four decades. Namely, we use here experimental recordings performed in one of the chief studies on intracellular pH regulation published by Roos and Boron in 1981 [24]. In this study (figure 5A of the original article), a *Helix* neuron was submitted to a 10 minutes 5%
CO₂ pulse. Its pH dropped from 7.35 to 6.85 and returned to normal after the pulse, with a noticeable overshoot. Details of the calculation and graphical results produced by our simulation are given in Materials and Figure S3A in File S1. Taken together, they show that only very minimal modifications of the constants of the system, well within differences found between different cell lines such as fibroblasts and neurons, have to be applied to converge to the resting intracellular pH measured in experimental conditions and that very satisfactory matches are obtained between calculated and experimental values for $\Pi_{\text{CO}_2}$ and intracellular pH.

**Main Outcome**

This study describes the first fully coupled and self-consistent mathematical system for intracellular pH regulation. For this, we constructed a minimal system that is uniquely based on the kinetic, electric and chemical equations describing the molecular processes pertaining to intracellular pH. This strategy is very different from classical heuristic methods used to model biological processes, that are mostly built on phenomenological equations deduced from fitted curves. It also avoids the introduction of ad hoc fluxes and/or constants to ensure the convergence of the numerical simulations with experimental data. Importantly, the present approach allows analytical processing. It shows, for the first time, that the dynamics of pH can be described by a reduced differential equation, and that steady-state intracellular pH values are in fact analytical solutions. Besides, despite the formal complexity provided by the large body of equations used here, the calculated numerical values of pH, ionic concentrations and membrane potential converge towards physiological values, with time evolutions that are very reminiscent of experimental behaviors. The last remarkable finding is the demonstration that any additional phenomenon that directly or indirectly impacts pH can be mathematically included without violating our model, provided that its equation is included without violating our model, provided that its equation is not ill defined. At this step, it is important to notice that here, we added to our model the effect of the CO₂ pulse.

**Results.**

Analytical expression of the pH dynamics, Natural Overshoot.

**Methods.**

**Parameters.**

Na/K-ATPase currents: voltage and sodium dependence.

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**Author Contributions**

Conceived and designed the experiments: YB MA. Performed the experiments: LG. Analyzed the data: LC YB MA. Contributed reagents/materials/analysis tools: LC YB MA. Wrote the paper: LC YB MA. Constructed the model: LC YB MA. Designed the software used in the analysis: YB.

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