Bi and Sn Doping Improved the Structural, Optical and Photovoltaic Properties of MAPbI$_3$-Based Perovskite Solar Cells

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Abstract: One of the most amazing photovoltaic technologies for the future is the organic–inorganic lead halide perovskite solar cell, which exhibits excellent power conversion efficiency (PCE) and can be produced using a straightforward solution technique. Toxic lead in perovskite can be replaced by non-toxic alkaline earth metal cations because they keep the charge balance in the material and some of them match the Goldschmidt rule’s tolerance factor. Therefore, thin films of MAPbI$_3$, 1% Bi and 0%, 0.5%, 1% and 1.5% Sn co-doped MAPbI$_3$ were deposited on FTO-glass substrates by sol-gel spin-coating technique. XRD confirmed the co-doping of Bi–Sn in MAPbI$_3$. The 1% Bi and 1% Sn co-doped film had a large grain size. The optical properties were calculated by UV-Vis spectroscopy. The 1% Bi and 1% Sn co-doped film had small $E_g$, which make it a good material for perovskite solar cells. These films were made into perovskite solar cells. The pure MAPbI$_3$ film-based solar cell had a current density ($J_{sc}$) of 9.71 MA-cm$^{-2}$, its open-circuit voltage ($V_{oc}$) was 1.18 V, its fill factor (FF) was 0.609 and its efficiency ($\eta$) was 6.98%. All of these parameters were improved by the co-doping of Bi–Sn. The cell made from a co-doped MAPbI$_3$ film with 1% Bi and 1% Sn had a high efficiency (10.03%).

Keywords: perovskite solar cells; Sn; Bi; MAPbI$_3$

1. Introduction

Energy consumption in every area of daily life is increasing and nanotechnology has revolutionized every aspect of daily life. In 1998, the rate of energy consumption was at 12.7 TW. In 2050, it will be at about 26.4 TW. This indicates that in 2050, over 46% extra energy would be needed compared to in 1998. However, the consumption of energy in 2021 is already very high compared to 1998. As a result, energy consumption in 2050 will be at about 58.7 TW [1]. To fulfill this high demand of energy, different traditional methods have been adopted, which are creating a high pollutant problem on Earth. Currently, the majority of energy is generated by nuclear and thermal power plants in different nations. Fossil fuels account for around 70% of all energy output. As a result, harmful pollutants, such as carbon oxides (CO$_2$), nitrogen oxides (NO$_x$), sulphur dioxide and chlorofluorocarbon are created when fossil fuels are burned. Every year, the temperature rises by around 0.5 to 1.1 degrees Fahrenheit, causing global warming, floods and ozone layer depletion [2]. As a result, academics from all around the globe are looking into more environmentally friendly alternatives to help solve this issue [3].

To replace conventional energy sources, some nations have aggressively created green, renewable energy resources (e.g., wind, solar and tidal energy, etc.). Among other renewable energy sources, solar energy is one of the greatest and most simply adaptable energy
sources. A solar cell harnesses this energy source as electricity on a large scale [4]. Unfortunately, solar cell manufacturing costs are considerable, preventing large-scale manufacture. Over the past decades, worldwide research has been underway to develop high-efficiency, low-cost and easy-to-manufacture solar cells [5]. These cells use the photoelectric effect to convert light into electricity in a direct manner. Various kinds of solar cells, including perovskite, hybrid, polymer, dye-sensitized and organic solar cells have been developed during this period [6–10]. Researchers are particularly interested in perovskite solar cells (PSCs) because of their ease of production, cheap cost and high efficiency. The light-harvesting active layer is usually made of a perovskite compound, which is made from hybrid inorganic–organic lead or tin halide-based materials. Organic–inorganic perovskite materials’ enormous success may be attributed to a unique mix of characteristics, including their high coefficient of absorption [11], charge carrier mobilities [12], large lifetimes of exciton and diffusion lengths [13] and low exciton binding energy [14]. The most well-researched perovskite absorbers are methyl-ammonium-lead-trihalide (CH$_3$NH$_3$PbX$_3$). In the formula, the “X” is a halogen atom, such as bromine, chlorine or iodine. CH$_3$NH$_3$PbX$_3$ may have an optical band gap ranging from 1.5 to 2.3 eV depending on the halide concentration and preparation temperature [15,16].

Despite the fact that MAPbI$_3$ possesses extraordinary features, MAI-based PSCs are facing some key problems. For example, they are not as efficient as their competitors. Another problem in the field of perovskite material development is lead toxicity (Pb). Pb-based perovskites disintegrate under ambient conditions, generating hazardous chemicals, making their usage especially troublesome [17,18]. To improve the performance of these perovskite materials, many tactics have been applied. For example, the mixing of organic/inorganic cations not only improves stability, but is also used to increase the optical properties of the material [19]. Significant experimental and theoretical effort has been expended to explore metal cation alternatives for Pb in order to find innovative perovskites for solar applications. Therefore, it has been crucial to reduce the Pb by replacing it with suitable alternatives (e.g., Sr, Ge, Sn, etc.) in perovskites. The complete elimination of Pb from the perovskite lattice, on the other hand, significantly reduces the performance of a solar cell [20]. Furthermore, halogen doping has been found to improve the quality, crystallization and performance of perovskite thin-film devices [21]. It was also reported that the performance of PSCs can be increased by considering isovalent-based doping at random sites. As Sn has the least formation energy, the excellent optical absorption and favorable band gap of Sn-doped PSCs make them potential candidates for high-efficiency perovskite cells applications [22]. It was found that adding a Bi-dopant to CsPbI$_3$ increased the α-phase stability by 13.2%, opening the door for the production of complete inorganic perovskites for photonic and solar applications [23]. Smaith et al. reported that color change was mainly due to increasing defect states and a considerable improvement in the sub band gap-related density of the states [24]. Edoardo et al. reported that their results confirmed the induction of deep trap states by the Bi dopant and that Bi$^{3+}$ behaved like a deep electron trap. The observed absorption initiation redshift after the doping of Bi in MAPbI$_3$ is mostly due to the transition to Bi defect states, with the perovskite band gap remaining virtually unaltered and the PCE increasing [25].

MAPbI$_3$ is doped with Bi and Sn to increase the absorption and injection efficacy of MAPbI$_3$-based PSCs. Due to the comparable atomic radii of Bi, Pb and Sn, these metals are readily doped in MAPbI$_3$. Heterovalent ions may be an effective substitute for Pb$^{2+}$ due to their matching ion sizes and having an electric structure of 6s$^2$6p$^0$. Isoelectronic (6s$^2$) is a trivalence bismuth ion (Bi$^{3+}$) with substantially low toxicity compared to Pb$^{2+}$ and comparable ionic radii (118 pm for Pb$^{2+}$ and 102 pm for Bi$^{3+}$), demonstrating the chemical behaviour similarities between both cations [26]. The specific surface area size of the MAPbI$_3$ film is enhanced when Bi is doped, and the band gap decreasing has been found to be due to inhibited MAPbI$_3$ crystallites with the inclusion of the Bi dopant enhancing ability [26,27]. MAPI degradation is significantly slowed when a little amount of Bi$^{3+}$ is added. As a result, another transition element, Sn, is utilised as a dopant to achieve this
high degree of photoemission efficiency [28]. Because lead (Pb) falls into the periodic table’s group IV, other elements from the same group, such as tin (Sn) and germanium (Ge), are considered the best substitutes for the Pb site in lead-based halide perovskites. In the periodic table, Sn is immediately above Pb, and the ionic radius of Pb$^{2+}$ (119 pm) and Sn$^{2+}$ (110 pm) are practically identical. Fortunately, Sn has gained prominence for its position in the periodic table, as well as within group IV elements, as a result of the outstanding chemical and physical properties related to Pb [29]. As a result, Sn has become the most common alternative element to Pb. Despite these similarities, Sn-based halide perovskites are unstable under ambient conditions, so the complete substitution of Pb with Sn is currently being considered. Due to its inherent instability, Sn$^{2+}$ has the propensity to oxidize into Sn$^{4+}$ in the presence of moisture. As a result, partial substitution of Pb with Sn, rather than complete replacement, is regarded as a superior method for improving perovskite’s efficiency [30]. Sn improves the efficiency of solar cells and has exceptional resistance to damp, humid and salty situations. Sn also lowers MAPbI$_3$’s recombination rate while increasing the efficiency of the injection of electrons from perovskites to the conduction band (CB) [28]. When photons from sunshine enter into the cell, the electrons of the highest occupied molecular orbital (HOMO) of perovskite are activated and migrate to the lowest unoccupied molecular orbital (LUMO). Because MAPbI$_3$’s conduction band is lower than the perovskite’s LUMO level, electrons are transferred to the CB of MAPbI$_3$ [31].

Based on previous studies, a Bi doping concentration of more than 1% reduces the effectiveness of PSCs; thus, 1% Bi is employed as a dopant in the present investigation. Furthermore, 0, 0.5, 1 and 1.5% Sn is co-doped in a MAPbI$_3$ film to improve the efficiency of PSCs even further.

2. Experimentation

The solution of MAPbI$_3$ was made by mixing MAI and PbI$_2$ with a 3:1 molar ratio. Both materials were mixed in 2 mL DMF (Dimethyl formamide) and stirred overnight at room temperature for homogeneous mixing. For the preparation of the Bi–Sn co-doped MAPbI$_3$ solution, 1 wt.% BiCl$_3$, 0 wt.%, 0.5 wt., 1 wt. and 1.5 wt.% SnCl$_3$, MAI and PbI$_2$ were dissolved in DMF and stirred overnight for homogeneous mixing. In this way, five solutions were prepared, which were MAPbI$_3$, 1% Bi, and 0%, 0.5%, 1% and 1.5% Sn co-doped MAPbI$_3$. For the deposition of the films, five substrates (FTO-glass, Mark, Germany) were washed with de-ionized water (DIW), ethanol and acetone several times. Thin films of undoped and Bi–Sn co-doped MAPbI$_3$ were deposited on substrates (FTO glass) by using the spin-coating method. Almost eight drops of prepared solution were added to FTO substrates and spun for 10 s with the spin at 4000 rpm. The films were annealed for 10 min at 100 °C in an electric furnace after deposition.

The phase structure and particle size of Bi–Sn co-doped MAPbI$_3$ films were analyzed using an X-ray diffractometer (PANalytical X’Pert PRO MRD), which is a multipurpose X-ray diffractometer equipped with an Ni-filtered Cu Kα source (=1.5418) capable of both high-resolution and lower-resolution measurements over a wide range of thin film and powder samples. It was used at 40 kV and 40 mA. A Shimadzu UV-2101 spectrometer was utilized to optically characterize the films.

The same process as in our previous research was used to coat TiO$_2$ film on an FTO/glass substrate [32]. To prepare the perovskite solar cells (PSCs), the precursors used were spiro-OMeTAD 72 mg, TBP 36 L, stock solution 22 L, 520 mg of 1mL, chlorobenzene 1mL and distilled lithium bis-(trifluoromethyl sulfony) imide in acetonitrile, which were used to make the hole transport material (HTM). Finally, a thermally evaporated 80 nm Au layer was applied to the device. These solar cells had a 0.16 cm$^2$ active area. The geometry of the cell was glass/FTO/TiO$_2$/MAPbI$_3$/spiro-OMeTAD/Au.

The thicknesses of the TiO$_2$, perovskite, Spiro and gold contacts were about 40, 350, 100 and 80 nm, respectively, measured by a surface profilometer. A calibrated air mass (AM) 1.5 G with 100 mW/cm$^2$ intensity and a computer-controlled Keithley 2400 (Keithley Instruments, Inc., Cleveland, OH, USA) source meter were used in a solar simulator to
generate the J–V curve, which was used to test the properties of the solar cells (starting from 1.2 to 0 V and sweeping direction is reverse bias. The device was measured under 1 sun illumination 1.5 AM with a scan rate of 0.2 V/s).

3. Results and Discussion

3.1. XRD Analysis

XRD was used to examine the atomic arrangement of pure and co-doped MAPbI$_3$ films, as shown in Figure 1. The peaks of MAPbI$_3$ were observed at 2θ values of 14.09°, 19.97°, 28.05°, 31.09° and 43.07° corresponding to (110), (200), (220), (310) and (314), respectively (PDF # 00-021-1276). The tetragonal structure of MAPbI$_3$ was observed and confirmed. These results match well with the literature [33]. Peak values of 2θ at 14.18°, 20.0°, 28.04°, 31.95° and 43.03° with corresponding planes (110), (200), (220), (310) and (314) were observed in the 1% Bi doped MAPbI$_3$ film. The shifting of peaks and a high intensity have been confirmed in Bi doping [28]. The ions of Bi were alternated with the site of MAI or inserted interstitially in the lattice with greater exciton binding energy (B.E) (approximately 400 MeV) and anisotropic carrier mobility [34,35], as a result of which intensity was increased. At (0.5%) of Sn doped in 1% Bi doped MAPbI$_3$ films, the 2θ values were 14.24°, 20.0°, 28.6°, 31.96° and 43.3° with corresponding planes (110), (200), (220), (310) and (314), respectively. At (1%) of Sn doped in 1% Bi doped MAPbI$_3$ films, the 2θ values were 14.28°, 20.1°, 28.6°, 32.0° and 43.4° with corresponding planes (110), (200), (220), (310) and (314), respectively. At (1.5%) of Sn doped in 1% Bi doped MAPbI$_3$ films, the 2θ values were 14.29°, 20.0°, 28.6°, 31.99° and 43.4° with corresponding planes (110), (200), (220), (310) and (314), respectively. All of the peaks shifted towards the lower angle with doping and no peaks related to MASn, SnPb and SnI, etc., were observed. This confirmed the doping. The minor shifting of peaks towards the lower angle was due to the smaller ionic radius of Sn$^{2+}$ (110 pm) than Pb$^{2+}$ (119 pm) [28]. A higher intensity of the peaks was observed at 1% Bi and 1% Sn co-doped MAPbI$_3$. This showed that this film had high crystallinity and good electrical properties. In the 1.5% Bi and 1% Sn co-doped MAPbI$_3$ film, the intensity of the peaks was reduced due to which crystallinity was reduced.

![Figure 1. XRD spectra of pure and Bi–Sn co-doped MAPbI$_3$.](image-url)

Grain Size (D) and Dislocation Line Density ($\delta$)

Grain size (D) describes the diameter of each grain sediment. By contrast, dislocation line density ($\delta$) describes the defects or irregularities of the crystalline structure. Where dislocation line density is used to calculate the number of flaws in a sample, a lower $\delta$ value means that the film is of higher quality [36].

The following relationships are used to compute the D and $\delta$ of films [37,38].

\[
D = \frac{0.9\lambda}{\beta \cos \theta}
\]
Grain Size (D) and Dislocation Line Density (δ) 

Grain size (D) describes the diameter of each grain sediment. By contrast, dislocation line density (δ) describes the defects or irregularities of the crystalline structure. Where δ = \frac{1}{D^2}

(2)

Here, λ is the wavelength of Cu-Kα (1.5406 Å) radiation, β is full width at half maximum (FWHM) and θ is Bragg’s angle [36]. The D and δ of undoped, 1% Bi and 0.5%, 1% and 1.5% Sn co-doped MAPbI3 of the extremely sharp peak (110) are shown in Figure 2. The D and δ of the pure MAPbI3 film were 44.5 nm and 0.505 × 10^{17} m^{-2}, respectively. When 1% Bi was doped into MAPbI3 film, D increased to 47.8 nm and δ decreased to 0.437 × 10^{17} m^{-2}. This is attributed to a better crystal formation. Furthermore, this indicated the appropriate lattice location of Bi atoms within the MAPbI3 film structure [26]. The D and δ of 1% Bi and 0.5% and 1% Sn co-doped MAPbI3 films were (60.2 nm and 0.276 × 10^{17} m^{-2}) and (80.03 nm and 0.156 × 10^{17} m^{-2}), respectively. The high “D” confirmed the improved crystal structure and reduced the flaws of the crystal. The reason for the high ‘D’ is that Sn is an enhancer rather than an inhibitor [28]. The D of 1.5% Sn doped into 1% Bi-MAPbI3 film was reduced slightly to 52.1 nm and δ increased to 0.368 × 10^{17} lines m^{-2}. This drop in D implies a reduction in the film’s crystal structure quality. This is due to the fact that there are extremely few substitutional sites accessible at this concentration. This supports the observations of Wang et al., who discovered that doping significant amounts of Sn into MAPbI3 films reduces grain size. Table 1 displays the values of D and δ.

![Graph showing Grain Size vs Dislocation Line Density](image)

**Figure 2.** δ and D of pure and Bi-Sn co-doped MAPbI3.

**Table 1.** D and δ of pure and Bi-Sn co-doped MAPbI3 films.

| Films                  | Grain Size(D) (nm) | Dislocation Line Density (DLD) (m^{-2}) | d-Spacing(A°) |
|------------------------|--------------------|-----------------------------------------|---------------|
| MAPbI3                 | 44.5               | 0.505 × 10^{17}                         | 6.31          |
| 1%Bi-MAPbI3            | 47.8               | 0.437 × 10^{17}                         | 6.30          |
| (1% Bi + 0.5% Sn)-MAPbI3| 60.2               | 0.276 × 10^{17}                         | 6.28          |
| (1% Bi + 1% Sn)-MAPbI3 | 80.3               | 0.156 × 10^{17}                         | 6.25          |
| (1% Bi + 1.5% Sn)-MAPbI3| 52.1               | 0.368 × 10^{17}                         | 6.24          |

3.2. Optical Properties

The optical measurements, such as band gap energy, refractive index, extinction coefficients and dielectric constants of pure, 1% Bi and 0.5%, 1% and 1.5% Sn co-doped MAPbI3 films were studied using Shimadzu UV-2101 UV-Vis spectroscopy (University of Lahore, Lahore, Pakistan).

3.2.1. Band Gap Energy (E_g)

To calculate the E_g of pure MAPbI3 and co-doped MAPbI3 films with Sn–Bi, a graph was plotted between (αhv)^{1/2} and energy of photon (hv) as shown in Figure 3. In this graph, “α” is the absorption coefficient, “ν” is the light frequency and “h” is the constant
of planks [33]. The energy of the band gap is calculated by following the equation of the MAPbI$_3$ film using Tauc’s relation (3) [39].

$$a\alpha h\nu = A (h\nu - E_g)^n$$ (3)

In the above equation, $A$ is a constant and $n$ is an exponent. The exponent depends on the type of electronic transition. The value of the exponent for the direct band gap is denoted as $n = 1/2$ and for the indirect band gap as $n = 2$ [40]. Similarly, the value of the exponent represented as $n = 3/2$ for direct forbidden, and the value for indirect forbidden transitions $n = 3$, are utilized [41].

Undoped MAPbI$_3$ film had 1.66 eV of $E_g$. The $E_g$ was reduced to 1.64 eV with the doping of 1% Bi into MAPbI$_3$. This is because of the higher electro negativity, stronger covalent bonding strength of Bi compared to Pb, and because more electronic states are contributed near the Fermi level by Bi defects [22]. At 0.5% and 1% Sn doping, the $E_g$ is reduced to 1.63 eV and 1.56 eV, respectively. Sn has previously reduced the inside traps of the band gap in MAPbI$_3$ [42]. At 1.5% Sn doping into the 1% Bi-MAPbI$_3$, the $E_g$ was increased to 1.60 eV. This is owing to the fact that at this doping dose, grain size decreases. The large amount of dopant could not substitute the Pb vacancy; thus, it spread on the surface, as a result of which the $E_g$ was increased [16].

3.2.2. Refractive Index

The bending and traveling of light rays from one medium to another medium is understood with the refractive index ($n$) [43]. It discusses the polarizability of ions into materials [44]. Therefore, it is an important parameter for solar cell application.

To determine “$n$”, Vandamme and Hervé presented the refractive index relationship shown below [45].

$$n = \sqrt{1 + \left(\frac{A}{E_g + B}\right)^2}$$ (4)

Constants $B = 3.4$ eV and $A = 13.6$ eV are used in Equation (6). Pure MAPbI$_3$ has a refractive index of 2.86. Co-doped MAPbI$_3$ films with 1% Bi + 0.5%, 1% and 1.5% Sn have “$n$” as 2.87, 2.88, 2.91 and 2.89, respectively. The value of “$n$” in all films is increased by up to 1% Bi and 1% Sn co-doping. At 1% Bi and 1.5% Sn co-doping, the value of “$n$” is decreased because $E_g$ is increased. The high value of “$n$” indicates that material is optically dense and the speed of light slows in the material [46–48]. Therefore, more electrons will emit due to light falling on the material. This will increase the current density of the solar cell and finally enhance its efficiency.
3.2.3. Extinction Coefficient (k)

k is calculated for pure and Bi–Sn co-doped MAPbI<sub>3</sub> films [49].

\[ n = k \Delta \chi^* \gamma \]  

(5)

This is the proposed relationship between electro negativity (\(\Delta \chi^*\)) and refractive index. Here, \(\gamma = -0.32\) is constant [44]. k is dependent on n; the behavior of k and n is relatively similar [50]. The value of k for MAPbI<sub>3</sub> and 1% Bi, 0%, 0.5%, 1% and 1.5% Sn co-doped MAPbI<sub>3</sub> is 2.22, 2.88, 2.23, 2.20 and 2.21, respectively. The variation in “k” is due to doping, which refers to the change in the density of the film. This variation in “k” refers to a red shift, which is due to inter-band electronic transition when the Bi–Sn are co-doped in MAPbI<sub>3</sub> film [51]. The variation in inter-band electronic transition reflects the change in band structures, which confirms the substitution of Sn and Bi in the place of Pb in MAPbI<sub>3</sub> structures.

3.2.4. Dielectric Constants (\(\varepsilon_r\), \(\varepsilon_i\))

The refractive index describes how much slower light travels through a material than a vacuum, whereas dielectric constants describe how much is muted in a substance relative to a vacuum in the electric field. Because light is an electromagnetic wave, two parameters are linked to each other [52–54]. Light is electromagnetic wave; it has electric and magnetic fields. Therefore, the dielectric constant has two parts. One is the real part (\(\varepsilon_r\)) which behaves like the refractive index, and other is the imaginary part (\(\varepsilon_i\)), which corresponds to the extinction coefficient [11]. Therefore, the dielectric constant (\(\varepsilon\)) is given by [55].

\[ \varepsilon = \varepsilon_r + i\varepsilon_i \]  

(6)

\(\varepsilon_i\) and \(\varepsilon_r\) are computed using the following relationships [56]:

\[ \varepsilon_r = n^2 - k^2 \]  

(7)

\[ \varepsilon_i = 2nk \]  

(8)

The absorption loss ability of a specific wavelength is related to the imaginary part [11]. The real parts vary due changing layer thickness, shape, chemical composition and material anisotropy, etc. [57]. \(\varepsilon_r\) is increased and \(\varepsilon_i\) is decreased due to doping. The real part of the dielectric constant that has high value is necessary to increase the efficiency of solar cells because they determine the coulomb interaction between holes and electron pairs and the distance or magnitude of charge carriers, as well as any fixed ionic charges in the lattice. The \(\varepsilon_i\) value was greater than \(\varepsilon_r\) in our optical examination, as shown in Table 2. This indicated a strong interaction between photons and electrons as a result of which a red shift appeared and \(E_g\) decreased [58].

| Films                  | \(E_g\) (eV) | \(\varepsilon_r\) | \(\varepsilon_i\) |
|------------------------|--------------|-------------------|-------------------|
| MAPbI<sub>3</sub>      | 1.66         | 3.308             | 12.69             |
| 1% Bi-MAPbI<sub>3</sub> | 1.64         | 3.3085            | 12.74             |
| (1% Bi and 0.5% Sn)-MAPbI<sub>3</sub> | 1.63 | 3.67              | 12.38             |
| (1% Bi and 1% Sn)-MAPbI<sub>3</sub> | 1.56 | 3.63              | 12.80             |
| (1%Bi and 1.5%Sn)- MAPbI<sub>3</sub> | 1.60 | 3.46              | 12.77             |

3.3. J–V Measurement

The derived model, which has an exponential dependence of the photocurrent on applied voltage, adequately explains the features of solar cells [59].

Figure 4 depicts the current density–voltage curves of the MAPbI<sub>3</sub> and 1% Bi, 0%, 0.5%, 1% and 1.5% Sn co-doped MAPbI<sub>3</sub>-based perovskite solar cells.
Solar cells parameters, such as short circuit current density ($J_{sc}$), fill factor (FF), open circuit voltage ($V_{oc}$) and power conversion efficiency ($\eta$%) are measured by the J–V curve and tabulated in Table 3. The below relationship is used to calculate the photovoltaic efficacy of solar cells [60].

$$\eta = \frac{FF \times V_{oc} \times J_{sc}}{Pin}$$  \hspace{1cm} (9)

**Table 3.** Solar cell parameters of pure and Bi–Sn co-doped MAPbI3 films.

| Solar Cells                          | $J_{sc}$ (mA/cm²) | FF    | $V_{oc}$ (V) | $\eta$% |
|-------------------------------------|-------------------|-------|--------------|---------|
| MAPbI3                              | 9.69              | 0.587 | 1.08         | 6.14 ± 2% |
| 1%Bi-MAPbI3                         | 10.41             | 0.577 | 1.079        | 6.48 ± 2% |
| (1% Bi and 0.5% Sn%) MAPbI3         | 11.53             | 0.586 | 1.085        | 7.33 ± 2% |
| (1% Bi and 1% Sn%) MAPbI3           | 12.9              | 0.591 | 1.08         | 8.83 ± 2% |
| (1% Bi and 1.5% Sn%) MAPbI3         | 12.09             | 0.629 | 1.09         | 8.23 ± 2% |

The input power of the incident light is denoted by $Pin$. The FF is determined in the following manner [61]:

$$FF = \frac{(I_{max} \times V_{max})}{(J_{sc} \times V_{oc})}$$  \hspace{1cm} (10)

In the above equation, the maximum current is $I_{max}$, while the maximum voltage is $V_{max}$. The cell fabricated with pure MAPbI3 has a $V_{oc}$ of 1.08 V, $J_{sc}$ of 9.69 mA cm$^{-2}$, FF of 0.587 and efficiency of 6.14%. When light falls on MAPbI3, an electron is excited and moves towards the CB of MAPbI3. The hole left behind is moved towards the valence band (VB) of MAPbI3, where it moves towards the hole transport layer. Similarly, the electrons in the CB of MAPbI3 move towards the CB of the electron transport layer (i.e., TiO$_2$). This electron moves towards FTO and then the external circuit. Therefore, $J_{sc}$ and $\eta$ are observed. If this excited electron moved towards the hole rather than TiO$_2$, then FF would decrease and no efficiency would be observed. When 1% Bi is doped in MAPbI3, all the photovoltaic parameters are seen to increase and reach the values $V_{oc} = 1.079$ V, $FF = 0.577$, $J_{sc} = 10.41$ mA cm$^{-2}$ and an efficiency of 6.48%. This efficiency is greater than that of the undoped MAPbI3. Bi has a +3 state and Pb has a +2 state. This provides extra electrons to the MAPbI3 as a result of which $J_{sc}$ is increased and efficiency is increased [62]. $J_{sc}$ is dependent on the excitation of electrons by the interaction of photons. The electron is transferred to the LUMO level of MAPbI3 by gaining energy from the photons. Then, this electron jumps to the conduction band (CB) of TiO$_2$ by the driving mechanism for electron injection [63]. Therefore, $J_{sc}$ depends on charge carriers generated by the active layer and the charge collection on the electrodes, demonstrating the necessity of sunlight absorption.
by a particular semiconductor for charge generation. The critical parameters for optimizing the J\(_{\text{SC}}\) of PSCs are: active semiconductor band gap, absorbance and carrier mobility [64]. In XRD and UV measurement, it was found that grain size increased and E\(_g\) decreased due to doping. The large grain size was a result of low resistivity; therefore, the mobility of electrons is high in doped film. Moreover, the high value of J\(_{\text{SC}}\) indicates that the cell has a high electron injection efficiency due to the fewer differences in CB levels in TiO\(_2\) and MAPbI\(_3\).

When 0.5% and 1% Sn was doped into 1% Bi-MAPbI\(_3\) film, efficiency improved to 7.33% and 8.83%, respectively. Free electrons are available in the outer shell of Sn. Therefore, an increase in the transport rate of electrons and a delay in charge recombination is produced in Bi–Sn co-doped MAPbI\(_3\)-based solar cells [65]. This is confirmed by increasing J\(_{\text{SC}}\) and FF in Bi–Sn co-doped MAPbI\(_3\)-based solar cells. The high values of FF result in a reduction of the recombination rate. The photo-excited electron-hole pairs are separated by a doping-induced electric field in each grain, which increases the recombine life span of electron-hole pairs [66]. In the co-doping of 1% Bi and 1.5% Sn-based solar cells, the J\(_{\text{SC}}\), FF and \(\eta\) decreased to 8.23%. The large amount of Sn could not fit on the substitutional states. Therefore, the efficiency decreased. The FF of constructed solar cells is low. As a result of the creation of ionic defects, which operate as trap-assisted recombination sites and contribute to J–V hysteresis, these cells have a long recombination process. Therefore, the efficiency of the cells doped at 1% Bi and 1.5% Sn MAPbI\(_3\) was lower than those doped at 1% Bi and 1% Sn MAPbI\(_3\). Solar cells based on MAPbI\(_3\)-co-doped with Bi–Sn displayed good stability for eight days when exposed to humidity and ambient air at room temperature. Furthermore, even after five days of storage in the air, the gadget displayed very little performance loss.

4. Conclusions

In conclusion, undoped and Bi–Sn co-doped MAPbI\(_3\) films were successfully prepared using a solgel method. The impact of co-doping on the structural, optical and photovoltaic properties of MAPbI\(_3\) was explored. Changes in intensity and peak positions due to doping were investigated by XRD. The grain size and E\(_g\) were found to decrease with doping. The 1% Bi and 1% Sn co-doped MAPbI\(_3\) film had a large grain size (80.3 nm) and a small value of E\(_g\) (1.56 eV). The solar cell of the MAPbI\(_3\)-based film had \(V_{\text{oc}}\) (1.18 V), J\(_{\text{SC}}\) (9.71 mA cm\(^{-2}\)), FF (0.609) and \(\eta\) (6.98%). The efficiency was increased by co-doping. The highest efficiency (8.83%) was observed by preparing the cell with a 1% Bi and 1% Sn co-doped MAPbI\(_3\) film. This high efficiency resulted from the low recombination rate.

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