Sb NQR as a microscopic probe in Te doped correlated semimetal FeSb$_2$ : emergence of electronic Griffith phase, magnetism and metallic behavior

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121,123Sb nuclear quadrupole resonance (NQR) was applied to Fe(Sb$_{1-x}$Te$_x$)$_2$ in the low doping regime ($x = 0, 0.01$ and 0.05) as a microscopic zero field probe to study the evolution of 3d magnetism and the emergence of metallic behavior. Whereas the NQR spectra itself reflects the degree of local disorder via the width of the individual NQR lines, the spin lattice relaxation rate (SLRR) 1/T$_1$(T) probes the fluctuations at the Sb - site. The fluctuations originate either from conduction electrons or from magnetic moments. In contrast to the semi metal FeSb$_2$ with a clear signature of the charge and spin gap formation in 1/T$_1$(T)$\sim$exp$/\Delta k_B T)$, the 1% Te doped system exhibits almost metallic conductivity and the SLRR nicely confirms that the gap is almost filled. A weak divergence of the SLRR coefficient 1/T$_1(T)$\sim T$^{-n}$ \sim T$^{-0.3}$ points towards the presence of electronic correlations towards low temperatures. This is supported by the electronic specific heat coefficient $\gamma = (C_{el}/T)$ showing a power law divergence $\gamma(T) \sim T^{-n} \sim (1/T_1(T))^{1/2} \sim T^{-n/2} \sim C_{el}/T$ which is expected in the renormalized Landau Fermi liquid theory for correlated electrons. In contrast to that the 5% Te doped sample exhibits a much larger divergence in the SLRR coefficient showing 1/T$_1(T)$\sim T$^{-0.72}$. According to the specific heat divergence a power law with $n = 2 m = 0.56$ is expected for the SLRR. This dissimilarity originates from admixed critical magnetic fluctuations in the vicinity of antiferromagnetic long range order with 1/T$_1(T)$\sim T$^{-3/4}$ behaviour. Furthermore Te-doped FeSb$_2$ as a disordered paramagnetic metal might be a platform for the electronic Griffith phase scenario. NQR evidences a substantial asymmetric broadening of the 121,123Sb NQR spectrum for the 5% sample. This has purely electronic origin in agreement with the electronic Griffith phase and stems probably from an enhanced Sb-Te bond polarization and electronic density shift towards the Te atom inside Sb-Te dumbbell.

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I. INTRODUCTION

Magnetic resonance is a very suitable microscopic tool for correlated matter at the verge of long range magnetic ordering and aims in particular to expose the real nature of the magnetic fluctuations (antiferromagnetic (afm) - versus ferromagnetic (fm)) by temperature- and field-scalings. Local moment 4f- and 3f- systems driven by RKKY - and Kondo interaction could be tuned towards order through the quantum critical point (QCP) by either pressure, substitution or magnetic field. Among 3d- magnets tunable quantum criticality could be found in itinerant systems like NbFe$_2$ and (Ta, V)Fe$_2$ but also in systems with more localized Fe moments like YFe$_2$Al$_2$ and YbFe$_2$Al$_2$. Here in contrast to the itinerant Fe systems there is strong evidence for the emergence of weak Kondo interaction among the localized Fe moments. Signatures of Kondo type of correlations are also found in some magnetic semimetals. FeSi$_2$, FeSb$_2$ and FeGa attracted great attention because of their non-magnetic ground state and their promising low temperature thermoelectric performance. Metallic behavior and Fe-based magnetism could be introduced by controlled substitutions on the Fe- or the framework- site. For example for Fe(Ga$_{1-x}$Ge$_x$)$_3$ Ga-NQR was performed to monitor the effect of Ge- substitution across the phase diagram and to probe the magnetic fluctuations at zero magnetic field via the spin lattice relaxation rate (SLRR) trough the QCP. In conclusion we found an absence of induced disorder, localized antiferromagnetic Kondo-like correlations at low doping levels and critical ferromagnetic fluctuations at the QCP. In contrast to that the Co substitution on the Fe site introduces antiferromagnetic correlations in (Fe$_{1-x}$Co$_x$)Ga$_3$ but with sizable induced disorder. Along this line we started to work on Sb NQR in Fe(Sb$_{1-x}$Te$_x$)$_2$ where in contrast to Fe(Ga$_{1-x}$Ge$_x$)$_3$ an electronic Griffith phase is predicted for the disordered paramagnetic metal at the verge of canted antiferromagnetism. Being a local probe at zero field NQR could capture both most relevant points a) the degree of disorder b) the onset of critical antiferromagnetic fluctuations at the verge of long range order. Furthermore NQR might be of use to disentangle the electronic from the magnetic Griffith phase. In the correlated electron metal picture the SLRR is strongly related to the specific heat coefficient $\gamma \sim \exp \left(\Delta k_B T\right)$ $\sim \gamma^2 \sim (C/T)^2$ (Korringa law). For weak itinerant metals the Moriya- and the Herz Millis theory captures many different cases and systems. Especially at the verge to long range magnetic order power laws for the SLRR are predicted $1/T_1(T) \sim T^{-3/4}$ (afm) $\sim T^{-4/3}$ (fm).
Frequently non Fermi liquid behavior (NFL) was found for many systems which frequently originates from by local disorder. Disorder induced NFL behavior is discussed for many 3d- and 4f- and 5f- systems. Here for itinerant 3d systems (and to some extend U- based 5f systems) the Griffith phase was established whereas for some more localized 4f Kondo systems the Kondo glass scenario was proposed to capture the effect of disorder on the bulk properties. So far detailed microscopic studies like NMR or μSR in quantum critical itinerant d electron semimetals on the local effect of doping and the evolution of an electronic Griffith phase are missing. Here we report the results of $^{121,123}$Sb NQR spectroscopy and nuclear spin-lattice relaxation (SLR) 1/$T_1$ experiments on the correlated semimetal FeSb$_2$ and the Te doped systems Fe(Sb$_{0.93}$Te$_{0.05}$)$_2$ and Fe(Sb$_{0.92}$Te$_{0.05}$)$_2$.

II. EXPERIMENTAL

Single crystals of Fe(Sb$_{1-x}$Te$_x$)$_2$ ($x = 0.01, 0.05$) were prepared as described in $^{24}$. For NQR measurements Fe(Sb$_{1-x}$Te$_x$)$_2$ single crystals which exhibit good metallic conductivity already at $x = 0.01$ were crushed into fine powder and mixed with paraffin. NQR experiments were performed using phase-coherent pulsed Tecmag-Apollo NMR spectrometer. $^{121,123}$Sb NQR spectra were measured using a frequency step point-by-point spin-echo technique at 4.2 K by integration of the spin-echo envelope in the time domain and averaging over scan accumulation number which depends on the sample. The $^{123}$Sb nuclear spin-lattice relaxation was measured using the saturation recovery method in the temperature range of 2.5 - 200 K. In addition, low temperature specific heat measurements were carried out on Fe(Sb$_{1-x}$Te$_x$)$_2$ ($x = 0.01, 0.05$) single crystals using the Quantum Design PPMS in the temperature range of 0.5 - 30 K.

III. RESULTS

A. $^{121,123}$Sb NQR spectra

Bulk measurements reported previously$^{24,25}$ provide only macroscopic evidence for the emergence of an electronic Griffith phase accompanied by NFL behavior in Te-doped FeSb$_2$. To obtain a microscopic insight into underlying physics of this system we performed $^{121,123}$Sb nuclear quadrupole resonance (NQR) spectroscopy study on the same Fe(Sb$_{1-x}$Te$_x$)$_2$ ($x = 0.01, 0.05$) samples. $^{121,123}$Sb NQR spectra measured at 4.2 K for both samples are presented in Fig.1 together with the spectrum of the undoped FeSb$_2$ at 10 K adopted from $^{24}$. As seen from this figure even a very small (1%) Te doping causes significant broadening of the Sb NQR lines. Moreover, $^{121}$Sb $v_1$ line (58.5 MHz; $|±1/2 \leftrightarrow ±3/2|$ transition) and $^{123}$Sb $v_2$ line (55.8 MHz; $|±1/2 \leftrightarrow ±3/2|$ transition) already start to overlap in the Fe(Sb$_{1-x}$Te$_x$)$_2$

\[ x = 0.01 \] compound. Further increase of Te doping ($x=0.05$) leads to complete overlapping of these two NQR lines and formation of two broad shoulders to the left from $^{123}$Sb $v_2$ NQR line. Similar asymmetric broadening with formation of a low frequency shoulder is exhibited by all other $^{121,123}$Sb NQR transition lines in the Fe(Sb$_{0.93}$Te$_{0.05}$)$_2$ sample (see Fig.1, upper panel).

![Fig. 1](color online). $^{121,123}$Sb spectra measured at 4.2 K in Fe(Sb$_{1-x}$Te$_x$)$_2$ compounds with $x = 0.01$ (lower panel) and 0.05 (upper panel). For comparison, the same Sb NQR lines for the undoped FeSb$_2$ measured at 10 K and retrieved from$^{24}$ are presented (lower panel). The intensities of all transitions except $v_2$ line (55.9 MHz; $|±3/2 \leftrightarrow ±5/2|$ transition) for the Fe(Sb$_{0.93}$Te$_{0.05}$)$_2$ sample are normalized on their maximum intensity values. Inset: $^{123}$Sb $v_1$ line (88.0 MHz; $|±5/2 \leftrightarrow ±7/2|$ transition) and $^{123}$Sb $v_2$ line (95.1 MHz; $|±5/2 \leftrightarrow ±7/2|$ transition) without normalization for the Fe(Sb$_{0.93}$Te$_{0.05}$)$_2$ sample. Solid lines are guides for eye.

The full width at half maximum (FWHM) for $^{123}$Sb $v_1$ line (44.85 MHz; $|±1/2 \leftrightarrow ±3/2|$ transition) amounts 0.45 MHz, 0.91 MHz and 3.63 MHz for the Te concentration $x = 0, 0.01$ and 0.05, respectively. In other words, only 5% of heterovalent doping of Te for Sb results in almost one order of magnitude Sb NQR line broadening which is rather substantial. For comparison, 5% of Co substitution for Fe in relative to FeSb$_2$ nonmagnetic Kondo-like semiconductor FeGa$_3$ causes increasing of $^{63}$Ga (1=3/2) NQR FWHM from 0.044 MHz to 0.18 MHz which is factor of 2 less than that in FeSb$_2$. Unfortunately, we were not able to estimate FWHM values for other $^{121,123}$Sb NQR lines due to line overlapping in the Fe(Sb$_{0.93}$Te$_{0.05}$)$_2$ sample. In order to extract quantitative information from experimental $^{121,123}$Sb NQR spectra we determined the line width at 80% level from maximum line intensity. The obtained values are listed in Table 1 demonstrating considerable increase in $^{121,123}$Sb NQR line width in FeSb$_2$ with Te doping.
TABLE I. Width of the $^{121,123}$Sb NQR transition lines in $Fe(Sb_{1-x}Te_x)_2$ samples determined at 80% from maximum line intensity.

| $Fe(Sb_{1-x}Te_x)_2$ | 121Sb | 123Sb |
|-----------------------|--------|--------|
| $I = 5/2$ | $\gamma/2\pi = 10.188$ MHz/T | $\gamma/2\pi = 5.517$ MHz/T |
| $Q = 0.36$ Barn | $Q = -0.49$ Barn |
| $\Delta \nu_1$ MHz | $\Delta \nu_2$ MHz | $\Delta \nu_3$ MHz | $\Delta \nu_2$ MHz | $\Delta \nu_3$ MHz |
| $x = 0$ | 0.11 | 0.14 | 0.19 | 0.07 | 0.06 |
| $x = 0.01$ | 0.34 | 0.38 | 0.49 | 0.13 | 0.30 |
| $x = 0.05$ | 1.39 | 1.17 | 1.53 | - | 1.48 |

B. $^{123}$Sb nuclear spin-lattice relaxation

To probe the effect of small Te doping on the dynamical properties of $FeSb_2$ system we performed $^{123}$Sb nuclear spin-lattice relaxation (SLR) measurements at $^{123}$Sb $\nu_2$ NQR line ($|\pm 3/2\rangle \leftrightarrow |\pm 5/2\rangle$ transition) as a function of temperature in the range of 2.5 - 200 K by means of saturation recovery method. We have selected this line to enable comparison with the SLR data for the undoped $FeSb_2$ semiconductor available only for $^{123}$Sb $\nu_2$ NQR line. Since only one NQR transition line was saturated, the $^{123}$Sb $I (1 = 7/2)$ magnetization recovery curves for $Fe(Sb_{1-x}Te_x)_2$ ($x = 0.01, 0.05$) samples were fitted by the sum of three stretched exponents:

$$M(\tau) = M_0 + \sum_{i=1}^{3} C_i [1 - \exp(-(2k_i W_i \tau)^n)]$$  \hspace{1cm} (1)

Here $M(\tau)$ is the spin-echo integrated intensity, $M_0$ is the remaining magnetization after the saturation comb (at $\tau \rightarrow 0$), $\tau$ is the delay time between the saturation comb and the spin-echo pulse sequence, $2W_0 = 1/T_1$ is the $^{123}$Sb nuclear spin-lattice relaxation rate, $C_i(\eta)$, $k_i(\eta)$ for $FeSb_2$ ($\eta = 0.43$) were taken from the numerical calculations and were assumed not affected by Te doping. The stretched exponent parameter $n$ was introduced in Eq. (1) to account for the structural disorder caused by Te doping. The examples of experimental recovery curves and their best fits to equation (1) obtained for $Fe(Sb_{1-x}Te_x)_2$ ($x = 0.01, 0.05$) samples at 4.2 K compared with that for the undoped $FeSb_2$ at 10 K (retrieved from Ref.22) are presented in Fig.2. As seen from this Figure, the approximation of the experimental recovery curves to equation (1) is rather good. While for binary $FeSb_2$ $n \equiv 1$, increasing of Te doping leads to significant decrease of the stretched exponent parameter: $n \approx 0.73(1)$ for $x = 0.01$ and $n \approx 0.64(2)$ for $x = 0.05$. This effect is a consequence of spatial distribution of $1/T_1$ values due to growing structural and magnetic disorder in $FeSb_2$ crystal lattice caused by Te substitution.

It is worth to mention two characteristic features seen from Fig.2. First, the remaining magnetization $M_0$ after the saturation comb (at $\tau \rightarrow 0$) is dramatically increasing with Te doping $x$: while the initial saturation is almost perfect in the undoped $FeSb_2$ ($M_0 \approx 0.04$), $M_0$ becomes $\approx 0.3$ for $x = 0.01$ and $M_0 \approx 0.54$ for $x = 0.05$. This effect reflects an extreme broadening and even overlapping of Sb NQR lines in $Fe(Sb_{1-x}Te_x)_2$ with increasing $x$ (Fig.1). The intense spin diffusion effectively hampers saturation process despite all our efforts to optimize the saturation comb and minimize the $M_0$ value. The second interesting feature of the experimental data shown in Fig.2 is significant visible shift of the recovery curves towards low $\tau$ values with increasing Te doping which indicates very fast increase of the $1/T_1$ values with increasing $x$. This effect also favors increasing of the remaining magnetization $M_0$, as have been observed for the undoped $FeSb_2$ sample with increasing temperature.

![FIG. 2. (color online). $^{123}$Sb magnetization recovery curves for the $\nu_2$ NQR line (quadrupole transition $|\pm 3/2\rangle \leftrightarrow |\pm 5/2\rangle$) in $Fe(Sb_{1-x}Te_x)_2$ ($x = 0.01, 0.05$) samples at 4.2 K and $FeSb_2$ at 10 K. The latter curve was adopted from Ref.22. Solid lines are the best fits to Eq.(2) with $n = 1, 0.73(1), 0.64(2)$ for $x = 0, 0.01, 0.05$, respectively.](image)
As has been shown in Refs. 24, 25, even extremely low Te doping of \( x = 0.001 \) leads to transition from semiconducting to metallic behavior so that at \( x = 0.01 \) one can expect Körrenga-like SLRN governed by conduction electrons. Indeed, for the \( FeSb_{0.99}Te_{0.01} \) sample \( 1/T_1 T(T) \) might be considered as almost temperature independent in the range of 2 - 70 K (Fig.3). Above 70 K \( 1/T_1 T \) in \( FeSb_{0.99}Te_{0.01} \) sample increases merging to that for the undoped \( FeSb_2 \).

![Fe(Sh\_x\_Te\_y)_2\_T\_Sb NQR at 3/2-5/2 transition](image)

**FIG. 3.** (color online). \( 1/T_1 T \) as a function of temperature for the \(^{123}\)Sb \( \nu_2 \) NQR line (\( | \pm 3/2 \rangle \leftrightarrow | \pm 5/2 \rangle \) ) in \( Fe(Sh_{1-x}Te_x)_{2} \) compounds \( (x = 0, 0.01 \text{ and } 0.05) \). Solid straight lines are the best linear fits according to formula: \( 1/T_1 T = a \times T^{-2(1+\lambda)} \) (see text).

For the \( FeSb_{0.95}Te_{0.05} \) sample where \( 1/T_1 T \) is one order of magnitude higher than for \( FeSb_{0.99}Te_{0.01} \) and a power-law divergence \( 1/T_1 T \sim T^{-0.72} \) \( (1/T_1 \sim T^{0.28}) \) (Fig.3) was found towards low temperatures.

C. Specific heat

In addition to the NQR spectroscopy data we performed low temperature specific heat measurements on the same \( Fe(Sh_{1-x}Te_x)_{2} \) \( (x = 0.01, 0.05) \) samples (Fig.4, upper panel). The data is in a rather good agreement with findings of Hu et al.\(^{26}\) on crystals from the same batch. Here power law divergences in \( \gamma(T) = C/T \) and \( \chi(T) \) are discussed in the framework of the disorder induced Griffith phase (GF) at the verge of magnetism.\(^{22,28}\) According to\(^{22,28}\) the low temperature divergence of specific heat in GF systems is described by power function \( C(T)/T = a \times T^{1+\lambda C} \) with \( \lambda C < 1 \). Then likewise\(^{22,28}\), the total low temperature behavior of specific heat in these compounds can be successfully fitted to the equation:

\[
C(T)/T = \alpha \times T^{-1+\lambda C} + b \times T^2 + c \times T^4 \tag{2}
\]

The second and third terms in Eq.2 describes harmonic and anharmonic contributions to specific heat, respectively.\(^{25}\) The obtained values of \( \lambda C \) for both samples \( (\lambda C = 0.88 \ (x = 0.01) \text{ and } x = 0.05 \) see tab 2) are in good agreement with that reported in Ref.\(^{25}\).

![C/T vs. T plot in Fe(Sb\_1\_x\_Te\_x\_2) compounds (x = 0.01 and x = 0.05) sample. Dashed and solid lines are the best fits to Eqn.(2) and (5), respectively.](image)

**FIG. 4.** (color online). Upper panel: \( C/T \) vs. \( T \) plot in \( Fe(Sh_{1-x}Te_x)_{2} \) compounds \( (x = 0.01 \text{ and } x = 0.05) \). Solid lines are the best fits to Eq. (1) (see text). Lower panel: Low temperature part of the \( C/T \) vs. \( T \) plot for the \( FeSb_{0.95}Te_{0.05} \) sample. Dashed and solid lines are the best fits to Eqn.(2) and (5), respectively.

| \( x \) | \( (\lambda C)^* \) | \( (\lambda C)_{exp} \) | \( (\lambda T_1)_{exp} \) |
|---|---|---|---|
| 0.01 | 0.91(7) | 0.88(4) | 0.90(6) |
| 0.05 | 0.72(3) | 0.70(4) | 0.64(4) |
IV. DISCUSSION

The origin of the observed low frequency shoulder at $^{121,123}\text{Sb}$ NQR lines can be understood as follows. With increasing of $\text{Te}$ content $x$ from 0.01 to 0.05 the number of hetero-dumbbells $\text{Sb-Te}$ also increases. These dumbbells are characterized by polarization of the $\text{Sb-Te}$ bond due to higher electronegativity of the $\text{Te}$ atom. Therefore, electronic density inside the $\text{Sb-Te}$ dumbbell is shifted towards $\text{Te}$ atom. As a consequence, a partial negative charge on $\text{Sb}$ atom is reduced causing decrease of EFG followed by decreasing of $\text{Sb}$ quadrupole frequency. At low $\text{Te}$ concentration ($x = 0.01$) this effect is not yet visible but 5% $\text{Te}$ seems to be enough for detection since the number of hetero-dumbbells $\text{Sb-Te}$ increases substantially and the left shoulder on $\text{Sb}$ NQR appears. In this simplified approach only the 1-st coordination sphere of $\text{Sb}$ is considered which might be visualized as a $\text{Sb-Sb}$ dumbbell with short interatomic distance ($\sim 2.8$ Å) surrounded by 6 $\text{Fe}$ atoms likewise strongly distorted octahedron (Fig.5). In 5% $\text{Te}$ substituted $\text{FeSb}_2$ sample one can expect appearance of hetero-dumbbell $\text{Sb-Te}$ in the 2-nd coordination sphere of homo-dumbbell $\text{Sb-Sb}$ which slightly reduce the charge on $\text{Sb}$. In conjunction with strong NQR line broadening caused by lattice disorder this explains why instead of separate peak to the left of main $\text{Sb}$ NQR line we observe just a left shoulder.

As seen from Table 1, for all samples the broadening of the $\nu_1$ line for $^{121}\text{Sb}$ isotope is higher than that for $^{123}\text{Sb}$ isotope in satisfactory accordance to the ratio of their quadrupole moments $^{121}Q/^{123}Q = 1.36$. This result provides an evidence of electronic quadrupole origin rather than magnetic origin of the $^{121,123}\text{Sb}$ NQR line broadening in $\text{Fe}(\text{Sb}_{1-x}\text{Te}_x)_{12}$ ($x = 0.05$) sample. This supports the concept of electronic Griffith phase (EGP) in $\text{Te}$ doped $\text{FeSb}_2$ and is in a strong contrast to the magnetic Griffith phase (MGP). Indeed in case of isolated magnetically ordered clusters characteristic for the magnetic Griffiths phase $\text{Sb}$ nuclei inside these clusters should exhibit strong hyperfine magnetic fields of about $0.1 \div 1$ T induced from electron spins localized on $\text{Fe}$. Than instead of pure NQR one should observe Zeeman perturbed NQR on $^{121,123}\text{Sb}$ nuclei with pronounced splitting (or, at least, strong broadening) of initial NQR transition lines which depends on value and orientation of internal magnetic field in respect to the main EFG axes and asymmetry parameter $\eta$ and is proportional to gyromagnetic ratio $\gamma$ (see, for instance $^{22}$). Since $^{121}\gamma/^{123}\gamma = 1.85$ broadening of $^{121}\text{Sb}$ NQR lines should be almost twice as for $^{123}\text{Sb}$ isotope. This is definitely not seen in our experimental $^{121,123}\text{Sb}$ NQR spectra in $\text{Fe}(\text{Sb}_{0.95}\text{Te}_{0.05})_2$ sample. In the upper limit of hyperfine magnetic field ($\sim 1$ T) induced on $\text{Sb}$ nuclei within ordered spin clusters of Griffiths phase one can even observe a “wipe-out” effect of disappearing of $\text{Sb}$ NQR lines originated from cluster volume due to their extreme Zeeman broadening. This effect should substantially reduce the total $\text{Sb}$ NQR intensity which was not observed in our experiment.

For a metal in the frame of the Landau Fermi liquid (LFI) the SLRR could be related to the specific heat via the density of states at the Fermi level which yields $1/T_1/1 = N(\varepsilon_F) \sim \gamma^2 \equiv (C/T)^2 \sim T^{2(1+\lambda)}$. Contrasting to that, if the metal is a weak itinerant magnet, the SLRR is more related to the low energy and $q$-averaged complex dynamic susceptibility $\chi(q,\omega)$ which yields $1/T_1/1 \sim \sum_q \chi(q,\omega)$. Here it matters if the correlations are $\sim \text{afm}$ (at $q = 0$) or $\sim \text{afm}$ (at $q \neq 0$). For $\text{afm}$ correlations the SLRR is frequently found to be proportional to the bulk susceptibility $1/T_1/1 \sim \chi \sim T^{-1+\lambda}$ for the MGP.

For the 1% sample the specific heat coefficient power law ($m \equiv 1 - \lambda = 0.12$) suggests a SLRR power law with $n = 2m = 0.24$ which is in rather good agreement with the experimental result ($n = 0.2$). For the 5% sample the specific heat coefficient power law ($m = 0.28$) suggests a SLRR power law with $n = 2m = 0.56$ which is much smaller than what is found by experiment ($n = 0.72$). This might point towards the fact that upon doping we have a crossover from more localized correlated metal to and $\text{afm}$ correlated itinerant metal at the verge of order. Here Moryia predicted a power law with $n = 3/4$ which is rather close to the experimental finding.

Nonetheless the specific heat coefficient enhancement factor at 2 K ($\gamma_{\text{SLRR}}/\gamma_{\text{SLRR}}$) is about 5 which suggests in the LFI theory an enhancement of the SLRR $(R = 1/T_1/1) (R_{\text{SLRR}}/R_{\text{LFI}}) \approx 25$ which is indeed experimentally confirmed by our spin lattice relaxation measurements.

Let us take a closer look to predictions for the SLRR in an Griffith phase. According to theoretical prediction for the magnetic Griffith phase the nuclear spin-lattice
the fluctuations with the spectrum based on scaling analysis of collective bosonic modes of
istic exponent value within the Tsvelik and Reizer model at low temperatures
\(\lambda\) is far from the experimental values \(x = 0.05\) based on specific heat data analysis. Although approximation of specific heat these models we revisited the low temperature specific
ature logarithmic divergence
function
more general description of NFL behavior predicts logarithmic rather than power divergence of specific heat at 
It is worth to note that the Tsvelik and Reizer model also predicts for the specific heat low temperature logarithmic divergence. However, \(-\ln T\) function diverges even slower than 
which suggests a electronic
The spin lattice relaxation results clearly show that the charge gap of the pure correlated semimetal \(FeSb_2\) is filled upon \(Te\) doping. In a first approximation based on the Landau Fermi liquid theory for correlated metals the low temperature divergence of the SLRR \(1/T_1T(T)\) could be scaled to the one of the specific heat coefficient \(\gamma(T)\). A very good agreement was found for the 1% sample whereas the 5% sample which is closer to the antiferromagnetic ordered phase shows a significant violation of the scaling. The power law coefficient \(n = 0.72\) is rather close to the one expected for antiferromagnetic criticality which is \(n_{afm} = 3/4 = 0.75\). Nonetheless the enhancement factor of 5 between the \(x = 0.01\) and \(x = 0.05\) in the specific heat yields an enhancement of 25 in the \(1/T_1T\) value which was experimentally verified. As for both samples the the specific heat divergence is in good agreement with Ref.\(^{28}\) which suggests a electronic Griffith phase the microscopic SLRR shows a remarkable lack of consistency with the magnetic Griffith phase predictions. Probably this is because of the \(q\)-averaging nature of the SLRR. Nonetheless strong evidence for antiferromagnetic critical fluctuations at zero field for the 5% \(Te\) doped \(FeSb_2\) sample is given. Antiferromagnetic criticality is a rare occurrence in \(Fe\)-based semimetals in general might provide a platform for further studies. Further more other local probes (like \(\mu\)SR) should be addressed to study Griffith phase systems.

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