STAR ClUSTERS IN PSEUDOBULGES OF SPIRAL GALAXIES*

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ABSTRACT

We present a study of the properties of the star-cluster systems around pseudobulges of late-type spiral galaxies using a sample of 11 galaxies with distances from 17 Mpc to 37 Mpc. Star clusters are identified from multiband Hubble Space Telescope ACS and WFPC2 imaging data by combining detections in three bands (F435W and F814W with ACS and F606W with WFPC2). The photometric data are then compared to population synthesis models to infer the masses and ages of the star clusters. Photometric errors and completeness are estimated by means of artificial source Monte Carlo simulations. Dust extinction is estimated by considering F160W NICMOS observations of the central regions of the galaxies, augmenting our wavelength coverage. In all galaxies we identify star clusters with a wide range of ages, from young (age \( \lesssim 8 \) Myr) blue clusters, with typical mass of \( 10^4 M_\odot \) to older (age \( > 100-250 \) Myr), more massive, red clusters. Some of the latter might likely evolve into objects similar to the Milky Way’s globular clusters. We compute the specific frequencies for the older clusters with respect to the galaxy and bulge luminosities. Specific frequencies relative to the galaxy light appear consistent with the globular cluster specific frequencies of early-type spirals. We compare the specific frequencies relative to the bulge light with the globular cluster specific frequencies of dwarf galaxies, which have a surface brightness profile that is similar to that of the pseudobulges in our sample. The specific frequencies we derive for our sample galaxies are higher than those of the dwarf galaxies, supporting an evolutionary scenario in which some of the dwarf galaxies might be the remnants of harassed late-type spiral galaxies that hosted a pseudobulge.

Key words: galaxies: bulges – galaxies: evolution – galaxies: spiral – galaxies: star clusters

Online-only material: color figures

1. INTRODUCTION

Over the last two decades, observational evidence has been accumulating that not all spiral galaxies possess a bulge resembling a small elliptical galaxy with a de Vaucouleurs \( R^{1/4} \) light profile. Bulges with a surface brightness profile that does not follow the de Vaucouleurs law also show disk-like, cold stellar kinematics (Kormendy 1993) and light profiles that are well fitted by an exponential or an intermediate Sérsic law (Andredakis et al. 1995). For these reasons, these bulges are known as pseudobulges or exponential bulges. Hubble Space Telescope (HST) imaging of pseudobulges also has revealed that they often contain central compact sources, most likely star clusters (Carollo et al. 1998) and that they have shallower nuclear slopes than traditional bulges with the same magnitude (Carollo & Stiavelli 1998). It is also intriguing that the nuclear cusp slopes of the exponential bulges are similar to those of dwarf elliptical galaxies with the same radii and luminosities. This similarity suggests that there might be an evolutionary link between pseudobulges and dwarf ellipticals, as if the presence of a disk did not affect the nuclear properties of pseudobulge spirals (Stiavelli et al. 2001).

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In order to study in detail the properties of pseudobulges and investigate possible formation scenarios, we isolated a sample of 11 pseudobulges from our previous HST surveys with WFPC2 and NICMOS and obtained for this sample additional HST/ACS imaging data. The relation between bulge properties and galaxy properties for these systems has been discussed already by Carollo et al. (2007). These authors found that the bulge properties correlate very well with those of the host galaxy, suggesting that evolutionary processes of the disk may be responsible for the formation of the pseudobulges. In this paper, we focus rather on the properties of the population of star clusters within the pseudobulges.

We have two main goals. The first is to characterize the star-cluster systems of these late-type pseudobulge hosts and to compare them with the growing body of literature focused on spirals with classical bulges (for a review see Brodie & Strader 2006). Under the assumption that globular clusters form in association with enhanced bursts of star formation (e.g., Beasley et al. 2002), the abundance of globular clusters is expected to depend on whether a bulge is formed “classically”; that is, with a single burst of star formation, or more slowly, through extended star formation associated with secular processes in the quiescent galaxy disk (Kormendy & Kennicutt 2004; Brodie & Strader 2006). In this latter scenario, a lower specific frequency would be expected.

The second goal is to continue to investigate the similarities between pseudobulges and dwarf ellipticals, first highlighted by
Stiavelli et al. (2001). A spiral galaxy can be stripped of its disk during close interactions with neighbors, especially effective in a group environment (Moore et al. 1996). This “galaxy harassment” scenario suggests evolution of late-type spirals with a pseudobulge into dwarf ellipticals. The stellar populations of low-mass galaxies in clusters are indeed consistent with some fraction having evolved in this manner (e.g., Conselice et al. 2003). Under this scenario, there also should be a correlation between the properties of the star-cluster systems of these two galaxy classes. Here we explore this connection by measuring the specific frequencies of globular clusters per bulge light in our sample and comparing it with those of dwarf ellipticals.

This paper is organized as follows. We briefly introduce the sample and discuss the data reduction in Section 2. Source detection and photometry are described in Section 3. Star clusters are identified in Section 4 based on a color and detection and photometry are described in Section 3. St a r classification, and galactic extinction are listed in Table 1. The effects are removed using the IRAF STSDAS task CRREJ and the remaining hot pixels were removed by interpolation. Finally, sky subtraction was performed by determining the sky values from the WF chips, in areas farthest from the nucleus.

The images in the F606W (V) filter were acquired between 1996 and 1997 by WFPC2 with the galaxy nucleus centered on the PC camera (the only WFPC2 chip considered in this study as its resolution is comparable with ACS/WFC), which has a scale of 0.046 pixel$^{-1}$ and a field of view of about 35"×35". The observations were carried out in fine lock with a nominal gain of 15 electrons/DN and a total exposure time per field of 600 s, split in two to allow the removal of cosmic rays. For these images the raw data were processed with the standard WFPC2 pipeline (CALWP2) in order to use the most recent reference frames for flat fielding, bias, and dark current subtractions; the cosmic rays were removed using the IRAF STSDAS task CRREJ and the remaining hot pixels were removed by interpolation. Finally, sky subtraction was performed by determining the sky values from the WF chips, in areas farthest from the nucleus.

The images in the F435W (B) and F814W (I) filters were acquired between 2002 and 2003. The observations were carried out with the ACS/WFC channel, whose pixel scale is 0.05 pixel$^{-1}$ for a total field of view of 3/4×3/4. Details of these observations can be found in Carollo et al. (2007). The standard ACS pipeline (CALACS) was used to perform the basic data reduction (flat-fielding, bias and dark subtractions, and removal of the overscan regions), then the IRAF STSDAS tasks ACSREJ and DRIZZLE were used, respectively, to achieve cosmic-ray rejection and to correct for geometric distortions. Finally, any remaining hot pixels were removed by interpolation and sky subtraction was carried out for all galaxies.

The images in the F160W (H) filter were acquired between 1997 and 1998 by NICMOS (Camera 2). The pixel scale is 0.075 pixel$^{-1}$ for a total field of view of about 19"×19". These observations were carried out in snapshot mode and split into multiple exposures (to allow better cosmic-ray rejection) for a total exposure time per field of 384 s for ESO 498G5, ESO 499G37, NGC 1345, NGC 1483, NGC 3259, and NGC 2758, and a total exposure time per field of 256 s for NGC 406, NGC 2082, NGC 3455, NGC 4980, and NGC 6384. These images were reduced through the standard pipeline software (CALNICA). The various NICMOS anomalies (e.g.,

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**Table 1**

Basic Parameters for the Sample Galaxies

| Name          | α(2000) (h m s) | δ(2000) (deg ′′) | R$_{tot}$ (mag) | D$_{tot}$ (Mpc) | M$_{Btot}$ (mag) | Type       | E(B − V) (mag) | R$_{search}$ (kpc) |
|---------------|----------------|------------------|----------------|----------------|----------------|-----------|---------------|------------------|
| ESO 498G5     | 09 24 41.10    | −25 05 33.0      | 13.96          | 37.7           | 32.88          | SX54P/SBbc| 0.107         | 3.254            |
| ESO 499G37    | 10 03 42.10    | −27 01 39.0      | 13.15          | 17.7           | 31.24          | SX57/SBc  | 0.075         | 1.529            |
| NGC 406       | 01 07 24.10    | −69 52 35.0      | 13.10          | 19.8           | 31.48          | SASS*Sc  | 0.024         | 1.708            |
| NGC 1345      | 03 39 31.70    | −17 46 42.0      | 14.02          | 19.1           | 31.41          | SB55P/SBc | 0.038         | 1.654            |
| NGC 1483      | 03 52 47.60    | −47 28 39.0      | 13.00          | 15.1           | 30.90          | SB4/Sb-Sc | 0.007         | 1.308            |
| NGC 2082      | 01 54 51.20    | −64 18 04.0      | 12.62          | 17.1           | 31.17          | SB3/SBb  | 0.058         | 1.481            |
| NGC 2758      | 09 05 30.80    | −19 02 38.0      | 13.46          | 31.3           | 32.48          | PSB4P/Sbc | 0.126         | 2.707            |
| NGC 3259      | 10 32 34.68    | +65 02 26.8      | 12.97          | 24.8           | 31.97          | SX4*Sbc  | 0.015         | 2.140            |
| NGC 3455      | 10 54 31.20    | +17 17 02.8      | 12.87          | 19.8           | 31.48          | PSX3/Sb  | 0.033         | 1.708            |
| NGC 4980      | 13 09 10.20    | −28 38 28.0      | 13.19          | 23.9           | 31.89          | SX1P/SBc  | 0.072         | 2.063            |
| NGC 6384      | 17 32 09.42    | +07 03 36.8      | 11.14          | 22.4           | 31.75          | SX4R/SBbc | 0.123         | 1.934            |

Notes. Right ascension (α), declination (δ), and total apparent blue magnitude (M$_{Btot}$) are from RC3 catalog (de Vaucouleurs et al. 1991). Distance (D) and distance modulus (dm) are from the online NASA/IPAC Extragalactic Database (NED). The absolute magnitude M$_{tot}$ is obtained from B$_{tot}$ and dm. Morphological classifications are from the RC3 (left) and from the UGC (right) catalogs. The color excess E(B − V) is taken from Schlegel et al. (1998). The last column lists the radius of the common ACS/WFC–WFPC2/PC field of view. $^a$ Using $H_0 = 73$ km s$^{-1}$ Mpc$^{-1}$. 

8 IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.
the pedestal anomaly) were corrected on a case-by-case basis. We direct the reader to Carollo et al. (2002) for further details.

3. PHOTOMETRY

3.1. Source Detections and Aperture Photometry

Our goal is to identify star-cluster candidates and to characterize their integrated photometric properties. Therefore, we aim at constructing a uniform catalog of sources, focusing on detections that lie in the common area of the F435W, F606W, and F814W images. The resulting search radius for each galaxy is given in the last column of Table 1. The NICMOS F160W images have a more limited field of view and we use them only to evaluate the possible impact of dust extinction (see Section 5.1).

Our study is limited to star clusters close to the nucleus of the host galaxy, where there are significant inhomogeneities in the image background, mainly due to dust lanes and gradients in the integrated-light profile. Therefore, we identify star-cluster candidates on a “variance-normalized” image, that allows improved detection of point-like sources in regions of rapidly varying background (Miller et al. 1997). The “variance-normalized” image is obtained by summing the F435W and F814W images to improve the signal-to-noise ratio (S/N) and then by dividing the result by a smoothed version of itself, realized with a $11 \times 11$ pixel median filter.

We use the IRAF task DAOFIND to search for star clusters in the images. This task, optimized for point-source detection, is adequate for our goal, due to the compact size (FWHM $\approx 1$ or 2 pixels) of the star clusters in our sample. Indeed, for our purposes, with our data, DAOFIND marginally outperforms extended-source detection packages such as SExtractor (Bertin & Arnouts 1996).

The coordinates of the sources detected on the “variance-normalized” image are then mapped into the local pixel coordinates of each image using the IRAF tasks XY2RD, RD2XY, and GEOMAP. We visually inspected all sources, with particular attention paid to those with FWHM $\approx 0\farcs25$ in the F435W band, corresponding to a half-light radius $R = 0\farcs15$ after deconvolution with the point-spread function (PSF), if the source has a Plummer surface brightness profile (see Section 3.2). The angular scale $R = 0\farcs15$ corresponds to $R \approx 13$ pc at 17 Mpc—the distance of our closest galaxy—and $R \approx 27$ pc at 38 Mpc—the distance of the farthest galaxy. These more extended sources have in general very low surface brightness and often appear to be artifacts and are thus excluded from the subsequent analysis. This cut is primarily motivated by the size distribution of galactic GCs, which have radii smaller than about 20 pc. In addition, all the sources in our final catalog have been visually inspected to remove bright foreground stars identified from their diffraction spikes.

We perform circular aperture photometry of all the detected sources using an aperture radius of 0\farcs15 for each image and estimating the background level in an annulus between 0\farcs5 and 1\arcsec in radius. We verify that this method yields a correct estimate of the background including the galaxy luminosity. We checked this by choosing the central part of the galaxies, where their luminosity gradient is steepest and comparing our photometry with the published photometry of Carollo et al. (1997, 1998), which took into account a detailed model of the galaxy light profile. In addition, the aperture correction and error estimates have been performed via Monte Carlo simulations on the actual images (see Section 3.2), therefore any residual bias is corrected to first order.

Table 2

| Filter            | Zeropoint (mag) | Threshold (mag) | Reddening (mag) |
|-------------------|-----------------|-----------------|-----------------|
| F435W (ACS/WFC)   | 25.77a          | 26.5            | 1.319           |
| F606W (WFPC2/PC)  | 22.08b          | 26              | 0.908           |
| F814W (ACS/WFC)   | 25.50a          | 26              | 0.586           |

Notes. Zeropoints used for calibrating the magnitudes, threshold magnitudes corresponding to S/N $\approx 5$, and reddening coefficients expressed as $A_{HST}/A_{VJC}$.

- Derived from Pavlovsky et al. (2006).
- Derived from Holtzman et al. (1995).

We consider only sources that in all bands have a magnitude greater than a chosen threshold magnitude, corresponding to an S/N greater than 5 (see Table 2). The photometric calibration is done by converting instrumental magnitudes to the VEGAMAG magnitude system by applying the photometric zeropoints listed in Table 2. All magnitudes are then corrected for Galactic extinction following Schlegel et al. (1998). Since the HST filters differ from the Johnson–Cousins bandpasses, reddening corrections are computed using the effective transmission curve for each HST filter, created with the IRAF task CALCBAND in the SYNPHOT package (see Table 2).

In this paper, we indicate with uppercase letters (e.g., V) the apparent magnitude in the HST filters, while we use symbols like $M_V$ for the corresponding absolute magnitude. Absolute magnitudes in Johnson–Cousins bands are instead indicated with symbols like $M_{VJC}$.

3.2. Aperture Correction and Uncertainty Estimates

Total magnitudes are determined by correcting for flux outside the measurement aperture. Aperture corrections, photometric uncertainties, and the completeness of our observations are estimated by means of Monte Carlo simulations. We add artificial star clusters of known luminosity and size to each image in random positions within the search area defined in Section 3.1, and we measure the retrieved flux with the same aperture photometry procedure adopted for the real sources. The luminosity profile of the simulated star clusters follows a projected Plummer law,

$$L(R) = \frac{L_{\text{tot}}}{\pi b^2} \left(1 + \frac{R^2}{b^2}\right)^{-2},$$

where $b$ is the assigned scale length and $L$ is the assigned total luminosity. This is then convolved with the PSF in the appropriate filter obtained with TinyTim (Krist & Hook 2001) before being added to the image.

The result for a representative photometric Monte Carlo simulation is shown in Figure 1. Note that the distribution of the output magnitude ($m_{\text{out}}$) at a fixed input magnitude ($m_{\text{in}}$) has a width much smaller than the difference $m_{\text{out}} - m_{\text{in}}$. We define the aperture correction as the difference between $m_{\text{in}}$ and the median of the distribution of the output magnitude $m_{\text{out}}$. We find that the aperture correction is almost independent of $m_{\text{in}}$, and depends mildly on the size of the artificial source, with variations of about 0.2 mag for clusters with observed half-light radii between 0\farcs05 and 0\farcs15. We therefore apply the aperture correction to the observed sources according to their apparent size, using linear interpolation across our grid of models.

The error on the photometry is similarly estimated from the Monte Carlo simulations as the variance of the output.
magnitudes of the artificial sources. The typical uncertainties are plotted in Figure 2 as a function of the magnitude for three galaxies representative of our sample, NGC 1483 ($D = 17.1$ Mpc), NGC 4980 ($D = 23.9$ Mpc), and NGC 2758 ($D = 31.3$ Mpc). Note that the photometric error so obtained is larger than the formal statistical photometric error reported by the IRAF aperture photometry task.

We neglect the effect of possible varying charge transfer efficiency (CTE) because the observations were carried out when the instruments were still young, thus the effect is not severe. Further, the background levels in the images are provided by the underlying galaxies, rather than just the sky, and are high enough to make the effect minimal (Holtzman et al. 1995).

The errors on the colors are also quantified using our Monte Carlo simulations. We place and recover artificial sources with a given input color ($M_B - M_I \equiv B - I = 1, 1.5, 2; M_V - M_I \equiv V - I = 0, 0.5, 1; M_B - M_V \equiv B - V = 0, 0.5, 1$) on the corresponding images. From the recovered photometry we measure the error on the color and correlate it with the errors in each band, under the assumption of a linear relation. For example, in the case of the $B - I$ color the error-relation reads

$$\sigma_{B-I}^2 = \sigma_B^2 + \sigma_I^2 - 2 \cdot r_{BI} \cdot \sigma_B \cdot \sigma_I,$$

where $r_{BI}$ is the linear correlation term. The typical uncertainties on the colors are plotted in Figure 3 as a function of the magnitude and the color itself for the three reference galaxies introduced above. Our analysis yields that the typical value of the correlation $r$ is only mildly dependent on the color itself, with variations from $r \approx 0.5$ to $r \approx 0.9$. We assume a reference value of $r \approx 0.7$ (the average value of the sample), which turns out to describe the color errors satisfactorily.

Finally, through the same Monte Carlo simulations we also determine the completeness of our observations. For this purpose, we define as “successfully recovered” a synthetic source whose output magnitude satisfies

$$m_{\text{in}} < m_{\text{out}} < m_{\text{in}} + m_{\text{corr}} + 2\sigma,$$

where $m_{\text{in}}$ is the assigned input magnitude, and $m_{\text{corr}}$ and $\sigma$ are respectively the aperture correction and the $1\sigma$ photometric error previously estimated. Note that (1) $m_{\text{corr}} \gg \sigma$ (for example, see Figure 1), thus a source which is recovered as brighter than its input luminosity has an error at several standard deviations
Figure 3. Photometric uncertainties in the color measurements relative to input magnitudes for representative input colors and galaxies in the sample, estimated from Monte Carlo simulations as in Figure 2. (A color version of this figure is available in the online journal.)

Figure 4. Completeness curves for our survey, determined via Monte Carlo simulations by placing and recovering artificial star-cluster sources with a Plummer surface brightness profile and half-light radius of 0.075′′, which corresponds to the mean size of the stellar clusters in our sample. (A color version of this figure is available in the online journal.)

Since we found that the completeness does not significantly depend on the radius of the objects, we show in Figure 4 only the completeness curves derived for the mean size of the stellar clusters (0.075′′) for three galaxies representative of our sample and we note that in all cases the incompleteness becomes severe only at \( V > 25.5 \).

4. IDENTIFICATION OF STAR CLUSTERS

To classify star clusters, we compare the colors of the sources in the final catalog with those expected from synthetic stellar populations. For this we use the spectral energy distribution of Bruzual & Charlot (2003) models based on the Padova (1994) tracks (Bertelli et al. 1994), a Salpeter (1955) initial mass function (IMF) with masses between 0.1 and 100 \( M_\odot \), a range of metallicities from 0.02 to 1 \( Z_\odot \) and ages from 1 Myr to 15 Gyr. In addition, we have also included self-consistently hydrogen and helium recombination lines as well as metal lines, as described in Oesch et al. (2007). The integrated model colors for the \( HST \) filters of the observations have been obtained by processing the synthetic spectra through the IRAF task CALCBAND within the SYNPLOT package. With the CALCBAND task we also include optional dust extinction (\( E(B-V) \leq 1 \) mag). For comparison, we also created Bruzual & Charlot (2003) models using a Chabrier (2003) IMF with the same mass cutoffs.

As expected, we verified that the resulting optical colors are identical in both cases (in fact the two IMFs are different only for stars below 1 \( M_\odot \)). The mass-to-light ratio is a factor \( \approx 1.8 \) smaller (with very low dependence on age and metallicity) when using a Chabrier IMF instead of a Salpeter IMF, due to the lower fraction of low-mass stars.

The Salpeter-IMF tracks in the \( V-I \) versus \( B-V \) color–color plane are shown in Figure 5 and compared with the observed sources. Based on these colors, most of the detected sources are consistent with being relatively young star clusters (age \( \approx 50 \) Myr). The data-model comparison for all galaxies, except NGC 6384, suggests that only a modest extinction is present (but see Section 5.1 for a more detailed discussion).

The spread of the observed sources in the color–color plane (standard deviation and central point of the distribution reported in Figure 5) is comparable with the photometric errors on the
Figure 5. Color–color diagrams ($B - V$ vs. $V - I$) for star clusters in each galaxy, showing a comparison of our data to Bruzual & Charlot (2003) models. Black dots are the data with the “old” (age $\geq 250$ Myr) and massive (mass $\geq 10^5 M_{\odot}$) clusters highlighted in bold. The large red dot corresponds to the nuclear star cluster of each galaxy. The models include emission lines and have solar metallicity, single stellar population, a Salpeter (1955) IMF and no reddening ($E(B - V) = 0$, blue line) or a 1 mag of reddening ($E(B - V) = 1$, green line). Along the theoretical tracks we mark the age of the synthetic population in Myr. The star cluster average colors with their 1σ dispersions are shown in red.

(A color version of this figure is available in the online journal.)
As in Figure 5 but for the color–magnitude diagram ($B - I$ vs. $M_I$) for each galaxy compared to the Bruzual & Charlot (2003) models with a Salpeter IMF. Here the tracks for the models are shown for a total model mass of $1 \, M_\odot$. Therefore, the difference in $M_I$ magnitude between an observed source and the model with a corresponding $B - I$ color gives the total mass of the observed star cluster. In addition, we show, as dashed lines, tracks without emission lines. The upper envelope of the data points is a selection effect given by the magnitude threshold of the survey.

(A color version of this figure is available in the online journal.)
colors (see Figure 3). Hence it is not possible to infer the star-cluster formation history from this plot.

A better diagnostic is provided by a color ($B - I$) versus magnitude ($M_I$) diagram (see Figure 6 where the $M_I$ magnitude of the model tracks is computed for a reference 1 $M_\odot$ mass). We can estimate the total mass of each observed source based on the difference in $M_I$ magnitude compared to the model with the same $B - I$ color. The $B - I$ color also highlights that up to 50% of the total detected sources are younger than 8 Myr (see Table 3) under the assumption of no dust (again, see Section 5.1 for a discussion of the possible impact of dust extinction). In all galaxies, the observed sources are clearly not distributed as an evolutionary sequence at constant mass. Instead the older the stellar population, the greater is its luminosity (and hence mass), because the mass-to-light ratio increases with age. This is consistent with a mortality scenario for star clusters in which only a small fraction of star clusters (usually the most massive) survives over time (Lamers et al. 2005, Gieles et al. 2007, Gieles & Bastian 2008).

Every galaxy in our sample has a number of young star clusters, with ages $\lesssim$ 8 Myr and masses of the order of $10^5 M_\odot$. These can be used to reconstruct the recent star formation history of the parent galaxies. Taken at face value, our results indicate an average recent star formation rate (SFR), in such clusters, of 0.1–0.01 $M_\odot$ yr$^{-1}$. Linking this SFR to the actual SFR is however challenging, both because we need to extrapolate our results to young clusters below the detection threshold and because the photometric uncertainties might introduce a Malquist bias difficult to quantify. Conservatively, we can consider our results a lower limit to the overall galaxy SFR.

In this paper, we focus on the properties of the oldest (age $\geq 250$ Myr) and most massive ($M \geq 10^5 M_\odot$) stellar clusters; for reasons discussed further below, we identify these as young candidate globular clusters. Our choice in age is an operational choice to avoid (1) the youngest systems, where stellar evolution of the most massive stars can influence the dynamical evolution of the system by their high mass-loss rates and (2) a knee in the luminosity–color relation (see Figure 6) that would lead to a degeneracy in the mass estimate for the sources. Once a metallicity for the stellar population is assumed, the age and mass selection translates into a single color and luminosity selection, respectively. If we assume a low metallicity ($Z = 0.02 Z_\odot$), the nuclear star cluster age turns out to be longer than the Hubble time, suggesting that these galaxies are more metal enriched (at least in their central regions studied here). For a metallicity range from 0.2 to 1 $Z_\odot$, the total number of sources selected as globular cluster candidates varies at the 15% level at most. In fact, the tracks at different metallicities are close to each other for objects older than 250 Myr (see Figure 7).

Therefore, in the following we adopt solar metallicity for the clusters. The age selection $\geq 250$ Myr then translates into a color selection ($B - I$) $\geq$ 0.73 mag, assuming no dust. The same color cut corresponds instead to a younger age selection (age $\geq 100$ Myr) if dust is present (see Section 5.1). Our mass selection ($M \geq 10^5 M_\odot$) translates into a cutoff luminosity that depends on the color of the source: all the sources whose magnitude is brighter than the corresponding 1 $M_\odot$ BC03 track, shifted by 12.5 mag, are accepted. Our selection has been carried out with the assumption of a Salpeter IMF. The luminosity cutoff we use refers instead to a mass $M \geq 5.6 \times 10^4 M_\odot$ if we were to assume a Chabrier IMF. We identify these older clusters as young candidate globular clusters since we calculate (Trenti et al. 2007) that they have more than a 50% probability of surviving tidal dissolution over a Hubble time, assuming they orbit within a point-like potential well. In addition, the age selection we apply (age $> 100$–250 Myr) is sufficiently long to rule out early dissolution by supernova feedback.

For each galaxy, the candidate globular cluster luminosity function (GCLF), i.e., the number of candidate globular clusters per unit of magnitude, is plotted in Figure 8. For all the galaxies, the GCLF (completeness corrected) is reasonably fitted by a Gaussian distribution, whose FWHM and peak magnitude are reported in Table 5. The residuals from the fit are somewhat asymmetric with an excess of faint sources. This might reflect an intrinsic skewness of the GCLF, which is however difficult to assert based on our data. Alternatively, it might be due to the presence of a small fraction of stellar contaminants with a luminosity distribution peaked at the faint end of the candidate globular cluster luminosity distribution (see Section 5.2). Note that dust extinction does not introduce skewness in the GCLF, but only changes its peak value unless the amount of dust present correlates with the luminosity of the star clusters.

### Table 3

| Name            | All | Age $\leq$ 8 Myr | Age $\geq$ 250 Myr |
|-----------------|-----|-----------------|-------------------|
| ESO 498G5       | 477 | 225             | 62                |
| ESO 499G37      | 434 | 230             | 53                |
| NGC 406         | 1058| 376             | 133               |
| NGC 1345        | 490 | 194             | 41                |
| NGC 1483        | 1208| 551             | 59                |
| NGC 2082        | 1401| 576             | 96                |
| NGC 2758        | 777 | 326             | 141               |
| NGC 3259        | 805 | 409             | 70                |
| NGC 3455        | 1317| 623             | 67                |
| NGC 4980        | 1070| 591             | 133               |
| NGC 6384        | 153 | 25              | 71                |

**Notes.** For each galaxy (first column) the total number of sources identified as star clusters is given in the second column, while the number of young (age $\leq$ 8 Myr) and old (age $\geq$ 250 Myr) star clusters is in the third and fourth column, respectively.
Figure 8. Solid lines: candidate GCLFs in the V band, completeness-corrected, with the best-fitting Gaussian superposed. Dotted lines show the GCLF without completeness correction. No extinction or contamination correction has been applied. The black dot corresponds to the nuclear star cluster of each galaxy. The histograms $\int_{MV} \frac{dN(M_V)}{dMV} dMV$ are normalized so that the total number of observed sources $N = \int_{MV} \frac{dN(M_V)}{dMV} dMV$. 
have three free parameters: the total stellar mass of the source, its population models with a variable amount of dust extinction. We field of view, we perform a minimum chi-square fit on the diagnostic to the complete sample. However, the more limited field of view of NICMOS does not allow us to apply this combination with the other bands we consider, it significantly impact of dust extinction. The F160W band is ideal, since in all the galaxies in our sample have NICMOS F160W photometry (F435W, F606W, F814W, and F160W) allowing for a variable amount of dust extinction with different extinction laws (Galactic: second column; LMC: third column; and SMC: fourth column). The last column reports the number of sources within the same field of view that are older than 250 Myr and more massive than 10^5 M⊙ when the fit is forced to have no dust extinction.

### Table 4

| Name       | (N > 100 Myr)_{G14−Est} | (N > 100 Myr)_{G15−Est} | (N > 100 Myr)_{G15−Est} | (N > 100 Myr)_{G15−Est} |
|------------|--------------------------|--------------------------|--------------------------|--------------------------|
| ESO 498G5  | 7                        | 3                        | 6                        | 5                        |
| ESO 499G37 | 2                        | 1                        | 1                        | 5                        |
| NGC 406    | 17                       | 14                       | 17                       | 23                       |
| NGC 1345   | 10                       | 8                        | 12                       | 18                       |
| NGC 1483   | 1                        | 1                        | 1                        | 10                       |
| NGC 2082   | 8                        | 3                        | 7                        | 18                       |
| NGC 2758   | 15                       | 11                       | 13                       | 20                       |
| NGC 3259   | 7                        | 5                        | 6                        | 14                       |
| NGC 3455   | 3                        | 3                        | 3                        | 7                        |
| NGC 4980   | 7                        | 5                        | 8                        | 19                       |
| NGC 6384   | 6                        | 5                        | 5                        | 6                        |

**Notes.** Number of star clusters more massive than 10^5 M⊙ and with age ≥100 Myr identified in the central region of the galaxies within the NICMOS F160W field of view. The number of sources has been obtained using a least chi-squared fit of the four-band photometry (F435W, F606W, F814W, and F160W) allowing for a variable amount of dust extinction with different extinction laws (Galactic: second column; LMC: third column; and SMC: fourth column). The last column reports the number of sources within the same field of view that are older than 250 Myr and more massive than 10^5 M⊙ when the fit is forced to have no dust extinction.

### Table 5

| Name       | M_{Vpeak} | FWHM     | N_{fit} | f_{cont} | f_{ext−bias} | N_{tot} |
|------------|-----------|----------|---------|----------|-------------|--------|
| ESO 498G5  | −9.32 ± 0.09 | 1.46 ± 0.14 | 67 ± 10 | 0.1 | < 0.05 | 57^{+10}_{−10} |
| ESO 499G37 | −7.69 ± 0.10 | 1.46 ± 0.15 | 53 ± 9  | 0.1 | 0.70  | 14^{+10}_{−13} |
| NGC 406    | −8.76 ± 0.09 | 2.15 ± 0.15 | 147 ± 15 | 0.1 | 0.30  | 92^{+20}_{−31} |
| NGC 1345   | −8.56 ± 0.14 | 1.80 ± 0.26 | 33 ± 6  | 0.1 | 0.45  | 16^{+7}_{−10}  |
| NGC 1483   | −6.87 ± 0.13 | 2.01 ± 0.22 | 61 ± 9  | 0.1 | 0.90  | 5^{+4}_{−4} |
| NGC 2082   | −8.16 ± 0.12 | 2.13 ± 0.20 | 84 ± 10 | 0.1 | 0.65  | 26^{+13}_{−20} |
| NGC 2758   | −8.75 ± 0.08 | 2.08 ± 0.13 | 171 ± 17 | 0.1 | 0.35  | 100^{+24}_{−39} |
| NGC 3259   | −8.45 ± 0.08 | 1.49 ± 0.14 | 72 ± 10 | 0.1 | 0.55  | 29^{+13}_{−16} |
| NGC 3455   | −7.91 ± 0.12 | 2.00 ± 0.19 | 69 ± 10 | 0.1 | 0.55  | 27^{+12}_{−16} |
| NGC 4980   | −8.12 ± 0.08 | 1.92 ± 0.14 | 139 ± 14 | 0.1 | 0.65  | 43^{+20}_{−22} |
| NGC 6384   | −8.34 ± 0.11 | 1.93 ± 0.19 | 76 ± 11 | 0.1 | 0.10  | 61^{+17}_{−12} |

**Notes.** Completeness-corrected GCLF fit for the galaxies in our sample (first column). Fitted peak M_{V} magnitude and FWHM are in the second and third columns, respectively. The fourth column gives the total number of candidate globular clusters as inferred from the fit (N_{fit}). The fifth column contains the estimated fraction of stellar contaminants in our sample (see Section 5.2). The sixth column gives the fractional bias in the number of old star clusters induced by dust extinction of their host galaxy (see Section 5.1). The last column reports our fiducial number N_{tot} of candidate globular clusters obtained, defined as N_{tot} = N_{fit} × (1 − f_{cont}) × (1 − f_{ext−bias}).

### 5. SYSTEMATIC UNCERTAINTIES

The results in Section 4 are derived under the assumption of no dust extinction present outside the Milky Way and of the absence of contamination in our sample. Here we discuss these two issues.

#### 5.1. Dust Extinction

All the galaxies in our sample have NICMOS F160W coverage in the central region, which we use to quantify the impact of dust extinction. The F160W band is ideal, since in combination with the other bands we consider, it significantly increases the wavelength baseline. However, the more limited field of view of NICMOS does not allow us to apply this diagnostic to the complete sample.

For the subsample of star cluster sources within the NICMOS field of view, we perform a minimum chi-square fit on the data in four photometric bands comparing them to single stellar population models with a variable amount of dust extinction. We have three free parameters: the total stellar mass of the source, its age, and its dust content. The fit is performed assuming different extinction laws within SYNPHOT: Galactic, LMC, and SMC. As expected, the best-fitting models with dust tend to be younger than their counterparts with no extinction. The estimation of the total mass of the sources is instead not strongly modified by allowing this additional degree of freedom.

We summarize the results of these fits in Table 4, where we give the number of sources older than 100 Myr and more massive than 10^5 M⊙, for different assumed extinction laws. These results are compared, in the same table, to those obtained assuming an age ≥ 250 Myr and no dust, and it may be seen that the estimates are within a factor of 2 of each other. The difference between an age of 100 and 250 Myr is not very significant within the context of infant star-cluster mortality: in fact, most mortality happens within the first 50 to 100 Myr and is connected to mass loss induced by the rapid stellar evolution of the most massive stars in the cluster (Parmentier 2009).

To take into account the effects of dust on the general sample, where NICMOS coverage is absent, we introduce a galaxy-dependent statistical correction f_{ext−bias} to the total number of
sources we define “massive” and “old” star clusters. $f_{\text{ext-bias}}$ (reported in the last column of Table 5) is computed within the NICMOS field of view as the fractional difference between the average number of sources identified in the presence of dust versus the number identified under the no-dust assumption. This correction for dust extinction is then applied when we use the number of old star clusters obtained from the no-dust scenario to derive their specific frequencies (see Section 6).

The correction $f_{\text{ext-bias}}$ is conservative. In fact, based on the likelihood ratio test, the statistical significance of dust extinction is above the 90% confidence level only for about one-third of the sources with NICMOS coverage. This means that for some of the sources the apparent detection of non-zero extinction might simply be a consequence of allowing an additional degree of freedom in the fit. Finally, the NICMOS field of view is located at the center of the galaxy, where the amount of extinction is maximal (e.g., see Holwerda et al. 2005). The star clusters outside the NICMOS field of view are thus expected to be less influenced by dust.

5.2. Sample Contamination

A second systematic source of uncertainty that needs to be evaluated is the presence of stellar contaminants in our sample. These could either be stars in the Milky Way too faint to be identified by the presence of diffraction spikes or super-giants in the host galaxy.

To quantify the impact of Milky Way interlopers, in each galaxy of our sample we selected a region of the ACS field as far away from the galaxy as possible, with an area comparable to that used to search for star clusters. On these external regions, we selected sources using the $B$ and $I$ bands according to the selection criteria we apply to old star clusters. The surface density of sources identified in the outskirts of the field of view is an upper limit to the number of Galactic contaminants. For all 11 galaxies, we obtained surface densities from 1% to 5% (average 2.5%) of that of old star clusters selected in our main search region, therefore we can rest assured that Galactic stars are a negligible source of contamination compared to the larger uncertainties related to the treatment of dust extinction.

To quantify the impact of bright stars in the host galaxy, we consider the models of Marigo et al. (2008) that provide luminosities for giant-branch stars. Even for very low metallicities ($Z = 10^{-4} Z_{\odot}$), the maximum luminosity of giant stars is $M_V > -7$; for higher metallicities, the peak luminosity decreases (Marigo et al. 2008). From our Figure 8, it can be immediately seen that at most a few percent of the sources we select as star clusters are fainter than $M_V = -7$. Brighter contaminants are possible if they are young, extremely massive stars. These rare hypergiants can reach $M_V \sim -10$ and exhibit a wide range of colors (e.g., see de Jager 1998 for a review). However, they are very short lived (a few million years) and thus not likely to be found outside the H II regions where they were born. Ideally, we would need images in a narrow filter centered around the H$\alpha$ emission line to identify these sources. As this is not available for our galaxies, we estimate their impact on the sample based on the measured SFR. Assuming a Salpeter IMF, we expect from a few to \(10\) stars with mass $M > 60 M_{\odot}$ (bright enough to reach $M_V < -7$, e.g., see Stothers \& Chin 1999) to enter in our selection for each galaxy. Combining these two sources of stellar contaminants, that is Galactic and extragalactic stars, we assume a fraction 10% of sample contamination (indicated as $f_{\text{cont}}$ in Table 5). While this might carry some uncertainty, its impact is certainly secondary compared to the one induced by extinction, except for ESO 498G5 and NGC 6384.

6. SPECIFIC FREQUENCY

For each galaxy, we calculate the specific frequency of the candidate globular clusters in our field of view, i.e., the older, more massive cluster population normalized to the host galaxy luminosity (see Table 6). This is defined as follows:

$$S_{\text{gal}} = N_{\text{tot}} \times 10^{0.4(M_{V_{\text{GC,gal}}}+15)},$$

where $M_{V_{\text{GC,gal}}}$ is the absolute magnitude of the galaxy within the search area in the Johnson–Cousins $V$ band and $N_{\text{tot}}$ (reported in Table 5) is the total number of candidate GCs, obtained by integrating the completeness-corrected GCLF and applying a correction $f_{\text{cont}} = 0.1$ for the estimated fraction of stellar contaminants and a galaxy-dependent correction $f_{\text{ext-bias}}$ to take into account dust extinction (see Section 5.2). The luminosity of the host galaxy within the star-cluster-search area in the $HST$-F606W filter is calculated through aperture photometry. The background level is estimated using the WF chips of WFPC2, which cover a larger region of the sky compared to the star-cluster-search area, limited within the PC chip. The F606W magnitude ($V$) is then converted to the Johnson–Cousins $M_{V_{\text{GC}}}$ absolute magnitude using

$$M_{V_{\text{GC}}} = V + 0.287 \cdot (V - I) + dm,$$

where $dm$ is the distance modulus, listed in Table 1 and the color $V - I$ is from Carollo et al. (2007), except for NGC 1345 (for this galaxy, the color is calculated from the photometry of the galaxy search area in both images). The coefficient 0.287 is obtained by fitting with a linear law the relation between $M_{V_{\text{GC}}}$ and $V - I$ for the Bruzual \& Charlot (2003) models of old star clusters (magnitudes computed with the IRAF task CALCBAND). Equation (5) agrees very well with the calibration given by Holtzman et al. (1995) in his Table 10 (maximum difference of 0.05 mag, in the range $-1 < V - I < 2$), but has two advantages. First, it is calibrated exactly on the color range we are interested in (see the caveats in Holtzman et al. (1995) on the limited range of validity for his conversion), and second our formula is calibrated on the class of sources in which we are interested (star clusters, i.e., adopting an IMF) rather than individual stars as in Holtzman et al. (1995). In computing the specific frequencies, we have excluded NGC 1483 (the closest

| Name       | $R_{\text{search}}$ (kpc) | $M_{V_{\text{GC,gal}}}$ (mag) | $S_{\text{gal}}$ | $M_{V_{\text{GC,bul}}}$ (mag) | $S_{\text{bul}}$ |
|------------|---------------------------|-------------------------------|-----------------|-------------------------------|-----------------|
| ESO 498G5  | 3.254                     | -19.04                        | 1.603           | -17.06                        | 8.605           |
| ESO 499G37 | 1.529                     | -17.08                        | 2.103           | -12.79                        | 11.003          |
| NGC 406    | 1.708                     | -18.01                        | 5.319           | -16.65                        | 20.623          |
| NGC 1345   | 1.654                     | -17.68                        | 1.406           | -17.57                        | 1.507           |
| NGC 2082   | 1.481                     | -18.20                        | 1.407           | -15.65                        | 14.070          |
| NGC 2758   | 2.707                     | -18.87                        | 2.806           | -17.20                        | 13.310          |
| NGC 3259   | 2.140                     | -18.55                        | 1.105           | -14.97                        | 30.133          |
| NGC 3455   | 1.708                     | -18.01                        | 1.708           | -13.66                        | 96.433          |
| NGC 4980   | 2.063                     | -18.29                        | 2.103           | -17.39                        | 5.00220          |
| NGC 6384   | 1.934                     | -19.56                        | 0.909           | -19.44                        | 1.00620          |

Notes. Candidate globular cluster specific frequencies with respect to the galaxy luminosity (fourth column) and to the bulge luminosity (sixth column). The relevant luminosities used here are given in the third and fifth columns and are expressed in the Johnson–Cousins system. The galaxy luminosity has been obtained by considering the light within a circular region of radius $R_{\text{search}}$ (second column), corresponding to the search area for the star clusters. The bulge luminosity is from Carollo et al. (2007).
galaxy in our sample) because its bulge is very extended spatially and it gives the dominant contribution to the light within the star-clusters search area.

The resulting values of $S_{N,gal}$ are listed in Table 6, where uncertainties on $S_N$ take into account the errors on the candidate GC count ($N_{gal}$), including those derived from dust extinction (see Section 5.1), but not those in the distance modulus as they are negligible compared to the other errors.

The definition of specific frequency adopted here is a generalization of the standard Harris & van den Bergh (1981) definition, which is based on the total luminosity of the galaxy and on the total number of GCs. We note though that we do not expect a large difference in the two definitions, because the PC camera of WFPC2 contains a large fraction of the total light of the galaxies. Indeed, the absolute visual magnitude $M_{V,gal}$ we measure in our sample differs by less than about 1 mag compared to the total $M_{V,gal}$ luminosity of the galaxy (estimated from the absolute blue magnitude reported in Table 1, assuming $B - V = 0.5$).

We also compute the specific frequency relative to the bulge luminosity ($S_{N,bul}$) defined as

$$S_{N,bul} = N_{tot} \times 10^{0.4(M_{V,bul} + 15)},$$

(6)

where again $M_{V,bul}$ is expressed in the Johnson–Cousins system, obtained as above from the bulge $V$ magnitude published by Carollo et al. (2007). The resulting frequencies of candidate globular clusters per bulge light are also shown in Table 6.

These frequencies for our sample of galaxies can be compared with published values for nucleated dwarf galaxies (Miller et al. 1998; see also Miller & Lotz 2007 and Peng et al. 2008), as well as for early- and late-type galaxies with classical bulges (see Brodie & Strader 2006). The $S_{N,gal}$ we measure in our late-type galaxies is qualitatively consistent with that of other spiral galaxies, but there is a moderate tendency toward higher $S_N$. This finding seems to be at odds with predictions from semianalytic models of globular cluster formation that link the specific frequency with the overall SFR of the galaxy (Brodie et al. 2002): based on their scenario we would have expected a lower $S_N$ for our sample compared to that of spiral galaxies with classical bulges. In fact, we expect a lower SFR in pseudobulges if they form from the inner disk. However, the Beasley et al. (2002) study aimed at reproducing the properties of GCs in elliptical galaxies, therefore the comparison is only indirect, based on the assumption that the bulge formation process is similar to that of small ellipticals. In addition, it did not include evolution of the GCs system of the simulated galaxies, but rather its properties were fixed and fine tuned to match the observed data at the time of birth.

Another possibility to reconcile with these theoretical expectations is that of disruption of our candidate globular clusters as they age. Most of the "old" star clusters in our sample appear in fact to be relatively young compared to Galactic globular clusters (ages from a few hundred Myr to 1 Gyr), so this suggests that tidal interactions with the parent galaxies will reduce the number of star clusters as they age. For example, Gnedin & Ostriker (1997) estimate that for the Milky Way, the Galactic globular cluster system has a half-life of the order of the Hubble time. This is consistent with detailed $N$-body simulations of the dynamics of star clusters in the presence of a tidal field (e.g., see Trenti et al. 2007 and references therein). Under a scenario of significant adult mortality for star clusters, the specific frequency for our sample will become lower than that of spirals with a classical bulge, as expected on the basis of theoretical modeling of star formation (SF): in the case of pseudobulges, with a likely extended but low-efficiency star formation history, fewer globular clusters are formed than in an equivalent burst of SF, such as that considered to create classical bulges (Kormendy & Kennicutt 2004).

Mean specific frequencies for nucleated dwarf galaxies ($\langle S_{N,dwarf} \rangle = 7.5$)—whose GCs are similar to candidate GCs in our sample—lie instead between our $\langle S_{N,gal} \rangle = 2.0$ and $\langle S_{N,bul} \rangle = 30$. In addition, the trend of frequencies of candidate globular clusters per bulge light versus luminosity is very similar to that of dwarf galaxies (see Figure 9). This is suggestive of a stripping scenario in which dwarf galaxies originate from the bulges of late-type spirals. One possibility is that galaxy–galaxy fly-by encounters (galaxy harassment: Moore et al. 1996) may strip away stars in the disk, preferentially leaving behind bulge members and the central star clusters, both more protected because they sit deeper in the potential well of the system. In fact, passive evolution of the bulges (aging them from 1 to 10 Gyr, to model dwarf galaxies that have older stellar populations than pseudobulges) and accounting for tidal dissolution and stripping of star clusters (50% tidal dissolution and 50% to 75% stripping) shift the candidate GC specific frequency versus luminosity relation to match that valid for dwarf ellipticals (see Figure 9).

### 7. NUCLEAR STAR CLUSTERS

A detailed study of the photometric properties of nuclear star clusters has been presented in Carollo et al. (1997, 1998, 2002), here we briefly discuss their inferred ages and dust content. A summary of the photometry for the central sources is reported in Table 7, while the resulting mass, age, and dust content from our maximum likelihood fit are in Table 8. The nuclear star clusters
Properties of the nuclear star clusters derived from the fit of photometric 
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2006; Rossa et al. 2006) and dynamically from spectroscopy 
E 
(NGC 1483 and NGC 2758) present a clear evidence that their 
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variable dust extinction as described in Section 5.1. Two galaxies 
This is not surprising as rejuvenation is observed in a large 
number of nuclear star clusters (Rossa et al. 2006; Walcher 
et al. 2001) or in situ formation (Milosavljevic 2004).

8. DISCUSSION AND CONCLUSION

In this paper, we studied the properties of (relatively) old star clusters in a sample of 11 late-type spiral galaxies selected for the presence of a pseudobulge (Carollo et al. 1997). Star clusters were classified based on their color and luminosity through comparison of population synthesis models with HST photometry in three bands (ACS F435W, F814W, and WFC2 F606W). By means of artificial source detections, we estimated a completeness of $\approx 50\%$ down to $V = 25.5$. The clusters in our sample present a wide range of ages and masses, from young blue clusters with ages of a few tens of Myr to an older, red population (age $> 100-250$ Myr). We focus on these older, red clusters and identify them as young globular cluster candidates. All these galaxies have massive nuclear star clusters with masses in the $10^7 - 10^8 M_\odot$ range. These sources are typically the brightest star clusters in their host galaxy. They have a relatively old stellar age ($\gtrsim 1$ Gyr) except for two cases where a younger and dusty stellar population is inferred. Overall, the properties of our sample appear consistent with both proposed formation scenarios for nuclear star clusters, namely merging of stellar clusters driven to the galaxy center by dynamical friction (Lotz et al. 2001) or in situ formation (Milosavljevic 2004).

The presence of young star clusters allows us to set a lower limit to the SFR in the galaxies, which turns out to be $0.1-0.01 M_\odot$ yr$^{-1}$. This continuous, low rate of star formation is consistent with the formation scenario for pseudobulges, postulated to arise out of secular processes in the disk (Kormendy & Kennicutt 2004), compared to a violent burst of SF needed to form a classical bulge.

Based on the Kormendy & Kennicutt (2004) discussion of the formation of pseudobulges, we would expect them to have a deficit of massive star clusters compared to spirals with classical bulges, but this is not what we find. In contrast, the specific frequencies (number of star clusters normalized
to the galaxy luminosity—Harris & van den Bergh 1981) for the old population are consistent, within our uncertainties, to published data for other spirals. This appears to be a solid result, especially since we considered only central star clusters (the WFCP2 high-resolution detector limits our area of search to the central 35”×35”), normalizing the specific frequency to the galaxy luminosity within the field of view: in general star cluster systems are more spatially extended than their host galaxy (e.g., see Djorgovski & Meylan 1994; Jordán et al. 2009), thus our specific frequencies are probably lower limits to the global specific frequency.

When the specific frequency is computed with respect to the bulge luminosity, we get even higher $S_N$. Interestingly, the specific frequency versus bulge magnitude $S_{NV,\text{bul}}(M_{V,\text{bul}})$ trend is similar to the one observed in dwarf ellipticals. Pseudobulges have photometric and kinematic properties very similar to dwarf ellipticals, thus it is suggestive that some dwarf ellipticals might be the result of evolution of spiral galaxies with pseudobulges: the galaxy disk might in fact be stripped during galaxy–galaxy interactions (galaxy harassment—Moore et al. 1996).

In this scenario star clusters are also stripped away, but still a sizable number might survive compared to the more spatially extended disk. At the same time the pseudobulge survives almost untouched by stripping, protected as it sits at the center of the galaxy potential well. While this scenario is overall appealing, detailed numerical simulations are needed for a proper validation. These will be presented in a follow-up paper.

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