Three-dimensional patterning of solid microstructures through laser reduction of colloidal graphene oxide in liquid-crystalline dispersions

Bohdan Senyuk, Natnael Behabtu, Angel Martinez, Taewoo Lee, Dmitri E. Tsentalovich, Gabriel Ceriotti, James M. Tour, Matteo Pasquali & Ivan I. Smalyukh

Graphene materials and structures have become an essential part of modern electronics and photovoltaics. However, despite many production methods, applications of graphene-based structures are hindered by high costs, lack of scalability and limitations in spatial patterning. Here we fabricate three-dimensional functional solid microstructures of reduced graphene oxide in a lyotropic nematic liquid crystal of graphene oxide flakes using a pulsed near-infrared laser. This reliable, scalable approach is mask-free, does not require special chemical reduction agents, and can be implemented at ambient conditions starting from aqueous graphene oxide flakes. Orientational ordering of graphene oxide flakes in self-assembled liquid-crystalline phases enables laser patterning of complex, three-dimensional reduced graphene oxide structures and colloidal particles, such as trefoil knots, with ‘frozen’ orientational order of flakes. These structures and particles are mechanically rigid and range from hundreds of nanometres to millimetres in size, as needed for applications in colloids, electronics, photonics and display technology.
Owing to unique physical properties\cite{1,2}, graphene-based materials and devices are shaping the future of photonics, electronics, optoelectronics and energy applications\cite{3–8}. For example, they promise to replace rare and brittle indium tin oxide transparent electrodes in the display industry\cite{9,10}. The most effective, low-cost and scalable approach to prepare graphene-based materials is the reduction of graphene oxide (GO)\cite{9,10}, which can be produced in abundance. Graphene oxide can be reduced by multiple methods\cite{5–26} including chemical reduction\cite{10}, photoreduction\cite{6–8,10–14,27} and even reduction mediated by biological microorganisms\cite{15}. Most of these methods require specific precursors and conditions, which make them expensive and unsuitable for large-scale production. Multiple lithographic methods\cite{5,16} have been developed for post-production of reduced graphene oxide (rGO) and its micropatterning into useful structures and assemblies, but they are time and resource consuming and limited by advances in mask production. Recently, reduction of GO and micropatterning was achieved using laser irradiation\cite{6–8,10–13,17–23,27}. However, its use so far has been limited to reduction and direct writing of patterns in dry two-dimensional (2D) solid films of GO.

Here we produce fully three-dimensional (3D) functional solid microstructures of rGO in an aqueous nematic liquid crystal (LC) of 2D GO flakes using pulsed near-infrared laser scanning and characterize them using nonlinear optical microscopy and other techniques. We show that photoluminescence of laser-reduced rGO is increased compared with pristine GO, which allows for using the same excitation beam for simultaneous reduction of GO into desired microstructures and their subsequent nonlinear photoluminescence imaging. Orientational order of LC phases promotes homogeneous internal structure of complex 3D rGO patterns. Our approach is reliable and scalable, does not require special conditions or chemical reducing agents and can be used in ambient conditions in high concentration aqueous dispersions of GO flakes. Because laser reduction with tightly focused femtosecond light is an intrinsically depth-resolved nonlinear process, it can produce layered, sandwiched, as well as fully 3D microstructures, including nonplanar surfaces such as those with the topology of trefoil torus knots, by programmable spatial steering of a tightly focused laser beam, without lithographic masks. Various 3D topologically nontrivial solid microstructures of rGO obtained with this approach can be used in electronic and optoelectronic devices and in basic research.

Results

Laser reduction of colloidal graphene oxide. Solid microstructures of rGO were produced in concentrated aqueous dispersions of pristine GO flakes\cite{28,29} (see Methods) sandwiched in between two clean, untreated glass substrates (Fig. 1a). Graphene oxide flakes (Fig. 2a) are hundreds of nanometres wide graphene sheets with a basal plane and edges decorated by oxygen-containing groups including carboxyls, hydroxyls, epoxides and others\cite{10}, which cause screened electrostatic repulsion that stabilizes the colloidal flakes against aggregation prompted by interlayer forces between flakes, making them both strongly hydrophilic and stable in colloidal dispersions. Assemblies of GO flakes at 0.25–1 wt% can be produced in abundance. Graphene oxide can be reduced by multiple methods\cite{5–26} including chemical reduction\cite{10}, photoreduction\cite{6–8,10–14,27} and even reduction mediated by biological microorganisms\cite{15}. Most of these methods require specific precursors and conditions, which make them expensive and unsuitable for large-scale production. Multiple lithographic methods\cite{5,16} have been developed for post-production of reduced graphene oxide (rGO) and its micropatterning into useful structures and assemblies, but they are time and resource consuming and limited by advances in mask production. Recently, reduction of GO and micropatterning was achieved using laser irradiation\cite{6–8,10–13,17–23,27}. However, its use so far has been limited to reduction and direct writing of patterns in dry two-dimensional (2D) solid films of GO. Laser reduction of colloidal graphene oxide. Solid microstructures of rGO were produced in concentrated aqueous dispersions of pristine GO flakes\cite{28,29} (see Methods) sandwiched in between two clean, untreated glass substrates (Fig. 1a). Graphene oxide flakes (Fig. 2a) are hundreds of nanometres wide graphene sheets with a basal plane and edges decorated by oxygen-containing groups including carboxyls, hydroxyls, epoxides and others\cite{10}, which cause screened electrostatic repulsion that stabilizes the colloidal flakes against aggregation prompted by interlayer forces between flakes, making them both strongly hydrophilic and stable in colloidal dispersions. Assemblies of GO flakes at 0.25–1 wt% can be produced in abundance. Graphene oxide can be reduced by multiple methods\cite{5–26} including chemical reduction\cite{10}, photoreduction\cite{6–8,10–14,27} and even reduction mediated by biological microorganisms\cite{15}. Most of these methods require specific precursors and conditions, which make them expensive and unsuitable for large-scale production. Multiple lithographic methods\cite{5,16} have been developed for post-production of reduced graphene oxide (rGO) and its micropatterning into useful structures and assemblies, but they are time and resource consuming and limited by advances in mask production. Recently, reduction of GO and micropatterning was achieved using laser irradiation\cite{6–8,10–13,17–23,27}. However, its use so far has been limited to reduction and direct writing of patterns in dry two-dimensional (2D) solid films of GO. Laser reduction of colloidal graphene oxide. Solid microstructures of rGO were produced in concentrated aqueous dispersions of pristine GO flakes\cite{28,29} (see Methods) sandwiched in between two clean, untreated glass substrates (Fig. 1a). Graphene oxide flakes (Fig. 2a) are hundreds of nanometres wide graphene sheets with a basal plane and edges decorated by oxygen-containing groups including carboxyls, hydroxyls, epoxides and others\cite{10}, which cause screened electrostatic repulsion that stabilizes the colloidal flakes against aggregation prompted by interlayer forces between flakes, making them both strongly hydrophilic and stable in colloidal dispersions. Assemblies of GO flakes at 0.25–1 wt% can be produced in abundance.

Figure 1 | Laser-induced reduction of GO flakes in aqueous dispersions. (a) Schematic diagram of the experiment. (b) Transmission-mode optical micrograph of an rGO plane (a dark square) in the aqueous dispersion of GO flakes obtained using 850 nm laser light; the inset shows the same sample area obtained using visible white light. (c) Photoluminescence texture of the rGO plane (a bright square) shown in b; the inset (6 × 6 µm²) shows an SEM image of a surface of the solid rGO plane microparticle retrieved from a discotic nematic LC sample. (d) Transmission and (e) photoluminescence along the yellow line, respectively, in b and c; magenta filled and empty triangles show intensities in corresponding points marked in b and c. (f) In-plane x-y and (g) cross-section y-z photoluminescence textures of three rGO planes produced in a thick (d ≈ 33 µm) sample of the aqueous dispersion of GO flakes; a vertical cross-section y-z texture in g was obtained along the line marked by yellow arrows and a letter g in f. Scale bars, 10 µm.
illuminated by the femtosecond laser light at a wavelength 850 nm (Supplementary Fig. 2a) and the inset shows the same texture when observed under white light. Broad-band nonlinear photoluminescence29 is detected from samples (Fig. 2e) under illumination by a femtosecond pulsed laser beam of low power (Methods and Supplementary Fig. 3) at 850 nm, which allows for nonlinear, intrinsically depth-resolved imaging of nematic textures33 in GO LCs29. On the other hand, illumination of a GO nematic with a pulsed laser beam of high laser energy density (fluence) results in a permanent physical change with properties drastically different from the pristine samples. We attribute this permanent change of illuminated area to a local reduction of GO flakes by a laser beam6–8,10–13,27.

Figure 1b shows a bright-field micrograph of a square area illuminated by a laser beam of high fluence, $E \approx 60-70$ mJ cm$^{-2}$. Its appearance is dark (Fig. 1b) and its transmittance $T$ is lower (Figs 1d, 3a) than through the unaffected bright area. At the same time, the laser-treated area produces significantly higher intensity of photoluminescence, $I_{PL}$ (Figs 1e, 2e, 3c), allowing for the simultaneous nonlinear photoluminescence imaging of a produced pattern (Fig. 1c). Owing to the intrinsic depth-resolved nature of this nonlinear optical process34,35, the laser reduction of aqueous GO is spatially localized in a small volume of about 300 nm in diameter defined by optical resolution33–35 and called a voxel. This 3D localization allows for voxel-by-voxel micropatterning of rGO not only in the plane of the sample (Fig. 1b,c,f) but also across the sample thickness (Fig. 1g), which sets our approach apart from previous techniques6–8,10–13,17–23,27 of photothermal or photo-induced reduction of GO and 2D micropatterning in thin dried films. Three planes of rGO, whose thickness is determined by optical resolution, were produced at different heights within a thick GO nematic dispersion sample (Fig. 1f) using a tightly focused laser beam. This is achieved without employing multi-layer lithography3. Programmable control of a scanning laser beam allows for precise voxel-by-voxel reduction of GO and micropatterning not only in-plane at different heights (Fig. 1f,e) but also continuous reduction across a thick layer of a GO nematic (Supplementary Fig. 4).

Characterization of rGO. To probe how a femtosecond pulsed laser beam at 850 nm modifies aqueous GO samples, we use ultraviolet–visible (Fig. 2e and Supplementary Fig. 2), Raman (Fig. 2f and Supplementary Fig. 5) and X-ray photoelectron (XPS; Fig. 2g,h) spectroscopies, as well as optical (Fig. 3a,b) and nonlinear photoluminescence (Fig. 3c,d) imaging (Methods). Raman spectra (Fig. 2f) collected from an untreated GO area and an rGO region (Fig. 1b) show pronounced peaks located at about 1,357 cm$^{-1}$ and 1,598 cm$^{-1}$ corresponding to a D band induced
by structural disorder and a G band associated with the vibration of sp²-bonded carbon atoms. The background fluorescence in a Raman signal is significantly quenched in the rGO areas (Supplementary Fig. 5), depending on the duration of irradiation or dwell time at constant fluence. There is no significant shift of D and G peaks or change in the ratio between their intensities, but the absolute intensity of the Raman signal is increased in the rGO areas (Fig. 2f). Compared with the GO areas, rGO regions show sharper D and G bands (Fig. 2f), which can be attributed to removal of oxygen-containing groups, as we additionally verify using XPS. Figure 2g,h shows XPS spectra, which clearly indicate the removal of oxygen-containing groups during the 850-nm laser irradiation of GO in water. XPS survey spectra (Fig. 2g) show signals of carbon and oxygen from both GO and rGO, however, the intensity of the oxygen peak O1s from an rGO area is significantly decreased, indicating that oxygen was partially removed. The high-resolution Cls XPS spectra (Fig. 2h) were deconvoluted into peaks corresponding to carbon in, respectively, C-C, C-OH, C=O and O=C-OH functional groups, which allows us to estimate the relative content of carbon not oxidized to oxygen as ~45.8% in GO and ~58% in the rGO area. The observed levels of reduction (Fig. 2g,h) are similar to those reported in previous studies of 2D laser reduction in solid thin films. In addition, measurements of resistance (Supplementary Fig. 6) of produced plain rGO stripe microstructures, which were performed using methodology similar to that developed for the study of photoreduction of GO in solid thin films, indicate that rGO microstructures are more conductive as compared with the electrically insulating GO.

Differences in physical properties between pristine and laser-reduced GO are also detected by optical spectroscopy (Fig. 2e) and optical bright-field and nonlinear photoluminescence microscopy (Fig. 3). The absorbance of rGO areas increases in the visible range and the onset of the absorbance peak is red-shifted (Supplementary Fig. 2b). The transmittance of rGO within the visible range is decreased (Fig. 3a) as they change appearance from transparent yellowish to opaque dark brown (Fig. 3b), controlled by the dwell time of the scanning beam. At the same time, the intensity of nonlinear photoluminescence from rGO areas increases (Figs 2e, 3c,d), with its maximum moving to shorter wavelengths, depending on the dwell time. The increase of photoluminescence from laser-reduced areas was also observed in dried films of giant GO flakes (Supplementary Fig. 7). GO can be laser-reduced even within just a single giant GO flake (Supplementary Fig. 7a). Interestingly, the transmittance, T, of rGO beyond a threshold irradiance continuously decreases almost linearly with irradiation time to a minimum (Fig. 3a,b), but the intensity of photoluminescence increases with irradiation time and reaches saturation after a rather short dwell time (Fig. 3c). This behaviour of T may provide insights into the physical origins of photoluminescence from graphene-based materials and can be related to increasing the fraction of sp²-hybridized carbon during reduction, which increases both absorbance and photoluminescence intensity. Optical characterization reveals a qualitative picture of the reduction process and formation of solid structures of rGO flakes. Pristine GO flakes (Fig. 2a) in the lyotropic nematic LC are oriented, on average, in the same direction and are kept apart by screened electrostatic repulsions between the charged GO surface, with the average spacing between the hydrophilic flakes determined by the volume fraction of GO in the dispersion (Fig. 2c). Laser-induced reduction largely removes oxygen-containing groups from the GO flakes (Fig. 2b,g,h) making them more hydrophobic and weakening the electrostatic repulsions (Fig. 2d). Consequently, hydrophobic rGO flakes overlap with each other and form a solid continuous structure (Fig. 2d and Supplementary Fig. 8c) defined by the 3D illumination pattern. Dependence of photoluminescence from rGO as well as white-light transmission on the irradiation time (Fig. 3) and microscopic observations (Fig. 1b,c) allow concluding that for the used parameters of the irradiation process the onset of reduction is not very sharp. Experiments also show that at high laser powers, E ≫ 100 mJ cm⁻², the process most probably can become more violent with rapid expulsion of water and oxygen and damaging integrity of a sample. We therefore constrain the irradiation parameters so that the onset of reduction and rGO flake aggregation into microstructures is smooth and continuous and dependent on duration of irradiation (Fig. 3). Importantly, the laser patterning is done through scanning of a tightly focused beam to reduce the structure voxel-by-voxel, without causing bubbles or disrupting orientational ordering of the flakes. Progressive reduction of GO flakes by the scanned beam focused to a volume < 1 μm³ and diffusion of flakes around the region of scanning allow us to avoid flows and non-homogeneity of the GO dispersion in the region of reduction. As the GO flakes within the micro-patterned structure collapse atop of each other due to lowering of electrostatic repulsion forces caused by reduction, water is displaced from inter-flake regions continuously during the scanning, so that the effect of these localized flows on GO flake ordering can be neglected. Solid structures of rGO are rigid and preserve their shape even under mechanical stress induced, for example, by the flow of LC during shear of confining glass plates (Supplementary Fig. 8a,b). The rGO structures produced in the bulk of the LC fluid are permanent colloidal objects and, at normal conditions, do not ‘dissolve’ in an aqueous GO dispersion. For example, such microstructures in the form of square-shaped sheets of rGO could be re-dispersed in water and then transferred onto solid substrates for scanning electron microscopy (SEM) imaging, an example of which is shown in the inset of Fig. 1c. Thus, this approach may allow for facile production of not only surface-attached micro- and nano-structures, but also complex-shaped colloids made of ribbons or tubes of rGO that could be then immersed into different fluid host media. SEM imaging of solid and rigid rGO microstructures (Supplementary Fig. 8c) retrieved from a sample after reduction show a fairly homogeneous surface morphology (the inset of Fig. 1c and Supplementary Fig. 9). The folded wrinkles and ordering of overlapped rGO flakes visible in the SEM images of the solid structures resemble smooth changes of orientation due to their on-average alignment with normals along the discotic nematic director in the original fluid nematic of GO, which was ‘frozen’ into a solid structure after flakes collapsed atop of each other due to weakening of electrostatic repulsions caused by laser reduction. Absence of visible deep cracks or pores suggests that the expulsion of water and oxygen during the laser-induced reduction, at discussed parameters, was not a rapid or violent process and did not introduce additional disorder or defects in the nematic molecular alignment field. Our findings indicate that, by leveraging the intrinsic ordering of GO flakes in a nematic phase, the shape anisotropy of ribbon-like cross-sections and the degree of reduction may potentially be controlled through varying the GO flakes concentration beyond the critical value of nematic formation.

Microstructures of rGO. Reduction of GO flakes in aqueous discotic nematic samples in a spatially confined nanosized spot allows for direct and precise micropatterning of various 3D solid structures of rGO (Figs 4, 5 and Supplementary Figs 4, 10) using computer-directed voxel-by-voxel scanning of a laser beam. The degree of reduction can be controlled by changing the fluence and irradiation/dwell time (Fig. 3 and Supplementary Fig. 10). Figure 4 shows different complex multi-level patterns of rGO.
produced in a bulk and at a surface of aqueous GO nematic and imaged by optical and photoluminescence microscopy. Using transmission-mode bright-field microscopy, one can see all the elements of micropatterns at the same time, but blurred depending on the location of the microscope’s focal plane relative to the features of micropatterns (Fig. 4a,d,l); an image is focused at a single level (Fig. 4a) of rGO ‘nanostairs’ going from the top to the bottom substrate, or it is focused on the different parts of the text (Fig. 4d,l) micropatterned on different levels in a GO nematic cell. At the same time, confocal photoluminescence microscopy allows the detection of the photoluminescence signal only from specific planes at different depths across the multi-level structures of rGO (see, for example, Fig. 4b,c,e,g), while other details outside of the laser scanned area are invisible (Fig. 4e,g and Supplementary Movie 1). This property of selective reduction can be used in high-capacitance all carbon6 information storage devices. Figure 4h–j shows an example structure that can be used as prototypes of thin-film 2D or 3D resistor elements. Intersecting multi-level or multi-layer structures (Fig. 4k–o and Supplementary Movie 2), where insulating GO is sandwiched in between the more conductive rGO layers (Supplementary Fig. 6), can be envisaged as prototypes of electronic elements for applications in electronics and optoelectronics, such as thin film capacitors6,18,20 and field transistors5. Importantly, our approach is suitable for direct mask-free production of not only in-plane or sandwiched structures and devices6,7,12,14,22 but also fully 3D interconnecting patterns, structures and nonplanar surfaces, which is one of the key advances as compared with, for example, in-plane structures presented in a pioneering study on photoreduction of GO in solid thin films11. To further demonstrate this capability, we produce an array of rGO trefoil knots (Fig. 5 and Supplementary Movie 3), which have topologically nontrivial 3D shapes described by a set of parametric equations as $r = L[2.1(\cos \phi - 2.25 \cos 2\phi), 2.1(\sin \phi + 2.25 \sin 2\phi), 6 \sin 3\phi)]$, where $L$ determines the outer size of the colloidal knot and the parameter $\phi \in [0, 2\pi]$. Bright-field
optical (Fig. 5a–c) and photoluminescence (Fig. 5e–g) images of the rGO trefoil knot are obtained when focusing at the top, middle and bottom of the knot, respectively; a 3D perspective view of the knot (Fig. 5h and Supplementary Movie 3) was reconstructed from experimental 3D data of photoluminescence imaging. Because the internal orientation of rGO flakes within the colloidal structures matches that of the surrounding GO flakes, the fabricated knotted particles differ from the polymerized ones34 in that they do not induce noticeable director distortions or topological defects (Fig. 5). This demonstrates that topological defects can be avoided when the easy axis (boundary conditions) for the director orientation on the surface of complex-shaped particles is such that it mostly matches that of the surrounding LC, even when the colloidal inclusions exhibit nontrivial surface topology of torus knots.

Discussion

We have demonstrated selective reduction of GO flakes in the bulk of an aqueous LC dispersion of GO flakes at ambient conditions using nonlinear excitation by a near-infrared femtosecond laser, which allows for mask-free micropatterning of 3D solid planar and nonplanar structures and surfaces of reduced GO. Besides its simplicity, this stable, reliable method does not require special conditions or chemical reduction agents and can be implemented at ambient conditions in aqueous dispersions of GO flakes. Advantages of micropatterning in a discotic nematic LC phase include the possibility of producing patterns having more homogeneous internal structure and the potential to attain ‘frozen’ orientational order of GO flakes, which can also be modified before reduction by applying an external field32 or using other stimuli. Our approach can be used for scalable production of graphene-based devices for photonics, electronics20, optoelectronics12, information22 and energy storage2 applications. Furthermore, solid colloidal particles of rGO produced here are enabling new explorations in other research areas such as, for example, for probing the interplay of topologies of surfaces, molecular fields and defects in soft condensed matter34. In this regard, the ribbon-like shape of thin rGO colloidal structures may allow for obtaining colloids with surfaces lacking orientability, such as Möbius strips39, which would be otherwise difficult to produce using techniques like two-photon polymerization34.

Methods

Preparation of aqueous dispersions of graphene oxide flakes. The improved, mostly single-layer GO flakes used in this work were synthesized by method28 modified for large-scale production and obtained as a dispersion in deionized (DI) water at a concentration of 2.5 g l−1 (0.25 wt%). The synthesis of GO flakes and their purification procedure was as following: 3 g of xGnP 5-µm-graphite nanoplatelets (XG Sciences, Inc.) and 200 ml of 9:1 H2SO4/H2O2 solution were added to a 1-L Pyrex beaker. The mixture was stirred to form 9 g of KMnO4. Slowly, 9 g of KMnO4 were stirred into the mixture, which turned a deep dark green colour. Precipitates were taken to avoid exceeding ~5 wt% of KMnO4 per addition and to assure that the change in colour from green to purple is complete before introducing more oxidant (this is important as solutions with more than 7% wt% of KMnO4 added to H2SO4 could explode upon heating). The mouth of the beaker was covered as much as possible using aluminum foil and the reaction was heated to 40 °C using a resistance-heated water bath. After 6 h, the mixture had thickened and turned from dark green to purple-pink. At this point, additional 9 g of KMnO4 were added and allowed to react for another 6 h. When the reaction was completed, the resulting slurry was quenched over 200 ml of ice-water containing 5 ml of 30% H2O2. The resulting bright yellow GO suspension was purified by a series of centrifugations and re-suspensions in a pure solvent. All centrifugations were performed at 4,000 r.p.m., for 90 min, and all re-suspensions were performed by shaking the centrifuged GO for ~4 h at 230 r.p.m. on a platform shaker using 200 ml of solvent. The product was centrifuged once as is, and centrifuged washed once in DI H2O once in 30% HCl, and twice more in DI H2O. To determine the concentration of the final solution, a 10-ml aliquot was filter-washed with 200 ml of three different solvents: methanol, acetone and diethyl ether. The final material was vacuum dried at 4 °C and room temperature overnight and weighed (25 mg).

Sample preparation. The aqueous dispersion of GO flakes was additionally tip-sonicated for 2 h at ~35 W of ultrasonic power using a Branson 250 sonifier (Micro-250 with a 125/375 W transducer) operating at 20 kHz (Branson Scientific) and equipped with a microtip of diameter 0.2 mm, which allowed for achieving monodisperse flakes of smaller size29. A homogeneous nematic LC phase (Supplementary Fig. 1) was obtained at higher concentrations of GO flakes in dispersions by controlled removal of excess water from dispersions subjected to centrifugation in a Sorvall Legend 14 centrifuge (Thermo Scientific) at 12,000 r.p.m. for ~10–15 min. LC dispersions of GO were sandwiched between two glass slides. Clean untreated glass substrates were used in samples for laser-induced reduction of GO flakes and optical and nonlinear microscopy observations. A gap between substrates was set using Mylar films (DuPont Teijin) of thickness 10-30 µm. Dispersions of GO flakes were sonicated in a Branson 8891 ultrasonic bath for ~5 min before assembling a cell. Evaporation of water during experiments was prevented by sealing samples with a ultraviolet curable glue (NOA 63, Norland Products, Inc.). To retrieve a produced rGO microstructure for SEM imaging, a sample was disassembled and the GO dispersion, together with the rGO microstructure, was released into a Petri dish with DI water. Then, after ‘washing’ it in DI water, the rGO microstructure with some water was soaked into a pipet tip and placed onto a silicon wafer (Supplementary Fig. 8c). Once the residual water evaporated, we performed SEM imaging.

Laser-induced reduction and optical characterization. The integrated 3D laser induced reduction and nonlinear photoluminescence imaging of GO aqueous samples was performed at room temperature using a multimodal nonlinear optical microscope33 coupled to an inverted Olympus microscope (Supplementary Fig. 3). A tunable (680–1,080 nm) Ti:sapphire oscillator (140 fs, 80 MHz, Chameleon Ultra II, Coherent) was used as an excitation source. The excitation beam was directed to the sample by a system of mirrors (DMs) and focused into the sample using the Olympus high numerical aperture (NA) oil immersion objective UPLSAPO 100/1.27 (O.11). The spatial 3D position of the excitation beam was controlled by a pair of Glan polarizer and half-wave plate to be ~1 mW in a sample to prevent the photo- and thermal damage. The excitation of GO flakes was performed at 850 nm (Supplementary Fig. 2a) and at 400–700 nm (Supplementary Fig. 2b) in a backward mode with a Hamamatsu photomultiplier tube H5784-20 (PMT1). The transmission-mode single wavelength bright-field image was collected in a forward mode with PMT2. The polarization of excitation could be varied using a half- or quarter-wave retardation plate mounted immediately before an objective. The same setup was utilized for a 3D reduction of aqueous GO flakes using a pulsed laser beam at 850 nm and laser fluence E <100 mJ cm−2. The dwell time, or scanning speed, was also controlled by a scanning unit FV300. Olympus Fluoview software was used for data acquisition and image reconstruction, and ImageJ software was used for data processing and analysis. Polarizing microscopy observations of samples in visible light and measurements of photoluminescence spectra were performed using the same IX-81 equipped with crossed polarizers, CCD camera (Flea, PointGrey) and a spectrometer USB2000 (OceanOptics) mounted onto a microscope. Raman spectra were measured using Renishaw inVia Raman microscope with excitation at 514.5 nm and low power to reduce unwanted heating or optical effects induced by a laser. XPS spectra were collected using a PHI Quantera XPS microprobe (Physical Electronics, Inc.) and SEM imaging was performed using a Zeiss Auriga FIB-SEM instrument.

References

1. Eda, G., Fanchini, G. & Chhowalla, M. Large-area ultrathin films of reduced graphene oxide as a transparent and flexible electronic material. Nat. Nanotechnol. 3, 270–274 (2008).
2. Bonaccorso, F., Sun, Z., Hasan, T. & Ferrari, A. C. Graphene photonics and optoelectronics. Nat. Photonics 4, 611–622 (2010).
3. Pang, S., Hernandez, Y., Feng, X. & Muller, K. Graphene as transparent electrode material for organic electronics. Adv. Mater. 23, 2779–2795 (2011).
4. Yong, V. & Tour, J. M. Theoretical efficiency of nanostructured graphene-based photovoltaics. Small 6, 313–318 (2009).
5. Li, F., Xue, M., Ma, X., Zhang, M. & Cao, T. Facile patterning of reduced graphene oxide film into microelectrode array for highly sensitive sensing. Anal. Chem. 83, 6426–6430 (2011).
6. Strong, V. et al. Patterning and electronic tuning of laser scribed graphene for flexible all-carbon devices. ACS Nano 6, 1495–1462 (2013).
7. Mukherjee, R., Thomas, A. V., Krishnamurthy, A. & Korakat, N. Photothermally reduced graphene as high-power anodes for lithium-ion batteries. ACS Nano 6, 7867–7878 (2012).
8. Li, H. & Buebeck, C. h. Photoreduction processes of graphene oxide and related applications. Macromol. Res. 21, 290–297 (2013).
9. Segal, M. Selling graphene by the ton. Nat. Nanotechnol. 4, 612–614 (2009).
10. Pei, S. & Cheng, H.-M. The reduction of graphene oxide. Carbon 50, 3211–3221 (2012).
11. Zhang, Y. et al. Direct imprinting of microcircuit patterns on graphene oxide films by femtosecond laser reduction. Nano Today 5, 15–20 (2010).
12. Petridis, C. et al. Post-fabrication, in situ laser reduction of graphene oxide devices. Appl. Phys. Lett. 102, 093115 (2013).
13. Huang, L. et al. Pulsed laser assisted reduction of graphene oxide. Carbon 49, 2431–2436 (2011).
14. Oraobana, E. et al. Holographic patterning of graphene-oxide films by light-driven reduction. Opt. Lett. 39, 4263–4266 (2014).
15. Tanizawa, Y. et al. Microorganism mediated synthesis of reduced graphene oxide films. J. Phys. Conf. Ser 352, 012011 (2012).
16. Hong, J.-Y. & Jang, J. Micropatterning of graphene sheets: recent advances in techniques and applications. J. Mater. Chem. 22, 8179–8191 (2012).
17. Sokolov, D. A., Shepperd, K. R. & Orlando, T. M. Formation of graphene features from direct laser-induced reduction of graphite oxide. J. Phys. Chem. Lett. 1, 2633–2636 (2010).
18. Gao, W. et al. Direct laser writing of micro-supercapacitors on hydrated graphite oxide films. Nat. Nanotechnol. 6, 496–500 (2011).
19. Tesh, H. F., Tao, Y., Tok, E. S., Ho, G. W. & Sow, C. H. Direct laser-enabled graphene oxide–reduced graphene oxide layered structures with micropatterning. J. Appl. Phys. 112, 064309 (2012).
20. El-Kady, M. F., Strong, V., Dubin, S. & Kaner, R. B. Laser scribing of high-performance and flexible graphene-based electrochemical capacitors. Science 335, 1326–1330 (2012).
21. Trusovas, R. et al. Reduction of graphite oxide to graphene with laser irradiation. Carbon 52, 574–582 (2013).
22. Li, X., Zhang, Q., Chen, X. & Gu, M. Giant refractive-index modulation by two-photon reduction of fluorescent graphene oxides for multimode optical recording. Sci. Rep. 3, 2891 (2013).
23. Abdelhayd, V. et al. Photothermal deoxygenation of graphite oxide with laser excitation in solution and graphene oxide increase in water temperature. J. Phys. Chem. Lett. 1, 2804–2809 (2010).
24. Chang, C. W., Hon, M. H. & Leu, I. C. Patterns of solution-processed graphene oxide produced by a transfer printing method. J. Electrochem. Soc. 159, H605–H609 (2012).
25. Zhuo, Q. et al. Large-scale synthesis of graphene by the reduction of graphene oxide at room temperature using metal nanoparticles as catalyst. Carbon N. Y. 52, 559–564 (2013).
26. Gengler, R. Y. N. et al. Revealing the ultrafast process behind the photoreduction of graphene oxide. Nat. Commun. 4, 2560 (2013).
27. Zhang, Y. et al. Photoreduction of graphene oxides: methods, properties, and applications. Adv. Optical Mater. 1, 10–28 (2014).
28. Marcano, D. C. et al. Improved synthesis of graphene oxide. ACS Nano 4, 4806–4814 (2010).
29. Senyuk, B. et al. Nonlinear photoluminescence imaging of isotropic and liquid crystalline dispersions of graphene oxide. ACS Nano 6, 8060–8066 (2012).
30. Kim, J. E. et al. Graphene oxide liquid crystals. Angew. Chem., Int. Ed. 50, 3043–3047 (2011).
31. Dan, B. et al. Liquid crystals of aqueous, giant graphene oxide flakes. Soft Matter 7, 11154–11159 (2011).
32. Shen, T.-Z., Hong, S.-H. & Song, J.-K. Electro-optical switching of graphene oxide liquid crystals with an extremely large Kerr coefficient. Nat. Mater. 13, 394–399 (2014).
33. Lee, T., Trivedi, R. P. & Smalyukh, I. I. Multimodal nonlinear optical polarizing microscopy of long-range molecular order in liquid crystals. Opt. Lett. 35, 3447–3449 (2010).
34. Martinez, A. et al. Mutually tangled colloidal knots and induced defect loops in nematic fields. Nat. Mater. 13, 258–263 (2014).
35. Gattass, R. F. & Mazur, E. Femtosecond laser micromachining in transparent materials. Nat. Photonics 2, 219–225 (2008).
36. Eda, G. et al. Blue photoluminescence from chemically derived graphene oxide. Adv. Mater. 22, 505–509 (2010).
37. Lui, C. H., Mak, K. F., Shan, J. & Heinz, T. F. Ultrafast photoluminescence from graphene. Phys. Rev. Lett. 105, 127404 (2010).
38. Shang, J. et al. The origin of fluorescence from graphene oxide. Sci. Rep. 2, 792 (2012).
39. Machon, T. & Alexander, G. P. Knots and nonorientable surfaces in chiral nematics. Proc. Natl Acad. Sci. USA 110, 14174–14179 (2013).

Acknowledgements

We thank F. Mirri, Q. Liu, C. Twombly, R. Trivedi, P. Ackerman, B. Dan, A. Lee, F. Vitale and C. Young for useful discussions. We are grateful to B. Chen (Rice University Shared Equipment Authority) and F. Mirri for kind assistance with XPS measurements. We also thank A. Sanders for his help with SEM imaging and acknowledge the use of the Precision Imaging Facility at NIST, Boulder for the SEM characterization reported in this work. This work was supported by the US Department of Energy, Office of Basic Energy Sciences, Division of Materials Sciences and Engineering under Award ER46921, contract DE-SC0010365 with the University of Colorado (B.S., A.M., T.L., I.I.S.), the Air Force Office of Scientific Research MURI Program, contract FA9550-12-1-0035 with the Rice University (J.M.T. and M.P.), the Air Force Office of Scientific Research FA9550-14-1-0111 (J.M.T.) and the Welch Foundation grant C-1668 (M.P.). B.S. also acknowledges support from the ICAM Branches Cost Sharing Fund and the Welch Foundation.

Author contributions

B.S., N.B., A.M., T.L. and D.E.T. performed experimental work. G.C. and J.M.T. provided graphene oxide materials. B.S., M.P. and I.I.S. analysed experimental results. I.I.S., M.P. and B.S. conceived and designed the project, and I.I.S. and B.S. directed the project. B.S. and I.I.S. wrote the manuscript with an input from all co-authors.

Additional information

Supplementary Information accompanies this paper at http://www.nature.com/naturecommunications

Competing financial interests: The authors declare no competing financial interests.

Reprints and permission information is available online at http://npg.nature.com/reprintsandpermissions/

How to cite this article: Senyuk, B. et al. Three-dimensional patterning of solid microstructures through laser reduction of colloidal graphene oxide in liquid-crystalline dispersions. Nat. Commun. 6:7157 doi: 10.1038/ncomms8157 (2015).