Muffin-like lanthanide complexes with an \(N_5O_2\)-donor macrocyclic ligand showing field-induced single-molecule magnet behaviour†

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Three mononuclear lanthanide complexes of a 2-pyridylmethyl pendant-armed 15-membered ligand (**(3,12-bis(2-pyridylmethyl)-3,12,18-triaza-6,9-dioxabicyclo-(12.3.1)octadeca-1,14,16-triene); \(L\)** with general formula \([Ln(L)(H_2O)(NO_3)](NO_3)_2\) (\(Ln = \text{Tb} (1), \text{Dy} (2), \text{and Er} (3)\) are reported. Based on X-ray diffraction analysis of 1 and 2, the central lanthanide atoms are nine-coordinated from the ligand \(L\) and one coordinated water molecule and one monodentate-bonded nitrate ligand. The coordination geometry of the \([LnN_5O_4]\) cores can be described as a muffin-like shape. Magnetic measurements revealed that all three compounds show field-induced single-molecule magnet behaviour, with estimated energy barriers \(U \approx 44–82\) K. The experimental study was complemented by CASSCF calculations showing a trend of an increasing first excited energy gap (Tb → Dy → Er) within the muffin-like geometry with the lowest magnetization tunnelling probability for the Dy\(^{\text{II}}\) complex 2.

Introduction

Lanthanide ions reveal many unique physical and photochemical properties, such as large magnetic moments and magnetic anisotropy as well as long-life luminescence with sharp emission lines.\(^1\) Therefore, their complexes have found many applications in luminescent sensors,\(^2\) MRI contrast agents,\(^3\) radionuclide therapy,\(^4,5\) and magnetically active liquid crystals.\(^6\) In the past few decades, more attention has been devoted to the magnetic properties of lanthanide complexes, because they were found as ideal candidates for single molecular magnets \(\text{SMMs}\)\(^7,8\) or single ion magnets \(\text{SIMs}\)\(^9,10\) due to their large magnetic moments and high magnetic anisotropy, and because of their potential applications in quantum computing,\(^11\) high-density memory storage devices and molecule spintronics\(^12\) and single molecule transistors.\(^13\)

Many complexes of lanthanides with pthalocyanines,\(^15\) Schiff bases,\(^16\) β-diketones,\(^17\) and polyoxometalates\(^18\) show a high effective spin-reversal energy barrier \(U_{\text{eff}}\) and blocking temperature \(T_B\), i.e. parameters usually characterizing the \(\text{SMM/SIM properties}\).\(^7,11\) An attractive alternative to the above-mentioned ligands may be macrocycles with a variable cavity size and donor atom set favouring the desired types/sizes of metal ions.

In the past, only a few complexes of lanthanides with 2-pyridylmethyl pendant armed polyaza- or polyoxa-aza macrocyclic ligands were prepared and studied (Scheme 1). The tetrakis(2-pyridylmethyl) derivative of cyclen \(L_1\) was studied in complexes with select lanthanides (La, Pr, Nd, Eu, Gd, Tb, Er, and Yb),\(^19,20\) which were nine-coordinated with approximately monocapped square-antiprismatic geometry (twisted for Pr, Nd and Eu). The complex cations were chiral and they crystallized as racemic compounds.\(^19,20\) An extended 18-membered hexaaza macrocyclic ligand with four 2-pyridylmethyl pendant arms \(L_2\) and its complexes with general formula \(\text{Ln}[(\text{Ln}(L_2))(\text{NO}_3)_2\text{H}_2\text{O}])_n\) where \(n = 2–4, \text{Ln} = \text{La}, \text{Ce}, \text{Pr}, \text{Gd}, \text{Tb}, \text{Er}\) and \(\text{Tm}\), were studied.\(^21\) Lanthanide atoms in \([\text{Ln}(L_2)]^\text{III}\) were ten-coordinated and showed a distorted bicapped square antiprismatic geometry. The 18-membered tetraoxa-diaza macrocyclic ligand \(L_3\) was studied only in its La\(^{\text{III}}\) complex and only in solution.\(^22\) Surprisingly, no magnetic or luminescence properties of the above-mentioned complexes have been investigated in detail to date. On the other hand, our attention was more devoted to rather unusual seven-coordinate lanthanide complexes, because they may represent a very promising geometry for optimal/effective preparation of lanthanide-based SMMs/SIMs. As it was published recently, the highest magnetization reversal barrier, even higher than \(1000\) K \((U_{\text{eff}} = 708 \text{ or } 1025\) K), was observed in seven-coordinate Dy\(^{\text{III}}\) complexes \([\text{Dy}(L_4)]\)\(^2\) where \(X = \text{Cl}^- \text{ or Br}^-\), with a pentagonal bipyramidal coordination sphere confirmed by X-ray analysis. Among these seven-
coordinate complexes, to the best of our knowledge, there are only three studies dealing with lanthanide complexes of 15-membered macromacrocyclic ligands (L5,24 and L6)25,26 and lanthanides (DyIII,24–26 ErIII,24–26) in which the central atom should adopt the pentagonal bipyramidal geometry, which unfortunately have not been confirmed by X-ray analysis yet. Nevertheless, the DyIII complexes [Dy(L5)Cl2]Cl·6H2O25,26 and [Dy(L6)Cl2]Cl·4H2O25,26 revealed large magnetic anisotropy which was reflected in their SMM/SIM behavior (Ueff = 7.75 and 23.7 K; T0 = 7.52 × 10−7 and 6.4 × 10−6 s, respectively).26 In this perspective, herein, we report the syntheses, X-ray crystal structures and magnetic properties of three mononuclear lanthanide complexes (TbIII, DyIII, and ErIII) with the 2-pyridylmethyl pendant-armed 15-membered macromacrocyclic ligand (L), which was recently synthesized and proved to provide seven-coordinate late first-row transition metal complexes, and even a Co(n) mononuclear SMM.27 In this work, we tested the coordination mode of L in lanthanide complexes and their magnetic properties.

Experimental section

Materials and methods

3,12-Bis(2-pyridylmethyl)-3,12,18-triaza-6,9-dioxabicyclo[12.3.1]octadeca-1,14,16-triene (L) was prepared according to the literature.27 Other chemicals were purchased from commercial sources and used without further purification. Elemental analysis (C, H, N) was performed on a Flash 2000 CHNO-S Analyzer (Thermo Scientific, Waltham, MA, USA). Infrared spectra (IR) were recorded on a Thermo Nicolet NEXUS 670 FT-IR spectrometer (Thermo Nicolet, Waltham, MA, USA) employing the ATR technique on a diamond plate in the range of 200–4000 cm−1. The mass spectra (MS) were collected on a LCQ Fleet Ion Trap mass spectrometer (Thermo Scientific, Waltham, MA, USA) equipped with an electrospray ion source and a three-dimensional ion-trap detector in the positive mode. The temperature dependent (T = 1.9–300 K, B = 0.1 T) and field dependent (B = 0–9 T, T = 2, 5, and 10 K) magnetization measurements were performed on a PPMS Dynacool (Quantum Design Inc., San Diego, CA, USA). Dynamic magnetic properties were studied by measuring ac susceptibility on a MPMS XL-7 SQUID magnetometer (Quantum Design Inc., San Diego, CA, USA). Powder XRD patterns were recorded with a MiniFlex600 (Rigaku) using Cu Kα radiation (λ = 1.5418 Å). Emission spectra were recorded in an acetonitrile/methanol mixture (V/V = 1:1) at room temperature using an AvaSpec-HS1024×121TTE spectrometer. The excitation source was a deuterium arc lamp.

Synthesis of [Tb(L)(NO3)(H2O)](NO3)2 (1)

To a stirred solution of L (71.0 mg, 0.16 mmol) in 10 mL of acetonitrile, Tb(NO3)3·5H2O (71.2 mg, 0.16 mmol) was added. The resulting solution was refluxed under stirring for 2 h. The obtained clear colourless solution was allowed to crystallize by diffusion of diethyl ether vapour at room temperature. After 3 days, colourless crystals (66 mg, yield 52%) were isolated by filtration on a glass frit, washed with cold diethyl ether (2 × 1 mL) and dried under vacuum over NaOH for the next 2 days. Anal. Calcd for C25H33N8O12Tb1: C, 37.70; H, 4.18; N, 14.07%. Found: C, 37.81; H, 4.22; N, 13.88%. MS, m/z (+): 434.28 [L + H]+ (Irel = 100%); 456.28 [L + Na]+ (Irel = 61%); 653.08 [(TbL(NO3)2)−2O + H]+ (Irel = 53%); 685.06 [(TbL(NO3)2)−2O + 2H]+ (Irel = 32%); 716.08 [TbL(NO3)3]+ (Irel = 39%). IR (ATR, cm−1): 831 (m), 950 (m), 1011 (m), 1032 (s), 1057 (m), 1082 (m), 1283 (s), 1296 (s), 1322 (s), 1370 (s), 1445 (m), 1605 (m), 2893 (s), 2909 (s), 2952 (s).

Synthesis of [Dy(L)(NO3)(H2O)](NO3)2 (2)

The synthesis follows the same procedure as described for 1 except that Dy(NO3)3·6H2O was used as the starting material instead of Tb(NO3)3·5H2O. Light pink crystals were isolated after 1 day (121 mg, yield 84%). Anal. Calcd for C25H33N8O12Dy1: C, 37.53; H, 4.19; N, 13.89%. Found: C, 37.64; H, 4.19; N, 13.89%. MS, m/z (+): 434.28 [L + H]+ (Irel = 100%); 456.28 [L + Na]+ (Irel = 67%); 658.08 [(DyL(NO3)2)−2O + H]+ (Irel = 16%); 690.07 [(DyL(NO3)2)−2O + 2H]+ (Irel = 11%); 721.08 [DyL(NO3)3]+ (Irel = 12%). IR (ATR, cm−1): 832 (m), 950 (m), 1011 (m), 1032 (s), 1057 (m), 1082 (m), 1283 (s), 1296 (s), 1323 (s), 1370 (s), 1446 (m), 1605 (m), 2894 (s), 2909 (s), 2954 (s).

Synthesis of [Er(L)(NO3)(H2O)](NO3)2 (3)

The synthesis follows the same procedure as described for 1 except that Er(NO3)3·5H2O was used as the starting material instead of Tb(NO3)3·5H2O. Light pink crystals were isolated after 1 day (111 mg, yield 86%)}.
306 triplets and 245 singlets for TbIII, 21 sextets, 224 quartets (RASSCF) method was employed in CASSCF calculations with model, except for O-attached hydrogens whose positions were determined by X-ray analysis were done with the MOLCAS 8.0 program package. The active space of the Ln3+ ion does not fit into the small macrocyclic cavity, and is twisted due to its high flexibility and because the large size of the Ln3+ ion does not fit into the small macrocyclic cavity, and by one oxygen atom from the water molecule (O3), and one oxygen atom from the nitrate ligand (O4) (Fig. 1A and S2A†).

**Results and discussion**

**Description of crystal structures**

Single-crystal X-ray diffraction analysis (for 1 and 2) and powder diffraction analysis (for 3) revealed that all the three compounds are isostructural (Fig. S1†), and crystallize in the tetragonal non-centrosymmetric space group. The crystallographic data and structure refinements for complexes 1 and 2 are given in Table 1, and the selected bond lengths and angles are listed in Table 2. The asymmetric units of 1 and 2 contain one [Ln(L)(H2O)(NO3)2]+ complex cation and two nitrate anions. The central LnIII atom is coordinated by the heptadentate ligand (L) with the N3O2 donor set, which is twisted due to its high flexibility and because the large size of the Ln3+ ion does not fit into the small macrocyclic cavity, and by one oxygen atom from the water molecule (O3), and one oxygen atom from the nitrate ligand (O4) (Fig. 1A and S2A†).

**Theoretical calculations**

The post-Hartree–Fock calculations performed on the [Ln(L)(NO3)(H2O)]([NO3]2) complexes 1–3 using the geometries experimentally determined by X-ray analysis were done with the MOLCAS 8.0 program package. The active space of the CASSCF calculations comprised of eight, nine and eleven electrons in seven metal-based orbitals for 1, 2, and 3, respectively. The Restricted Active Space Self-Consistent Field (RASSCF) method was employed in CASSCF calculations with the following numbers of multiplets: 7 septets, 140 quintets, 306 triplets and 245 singlets for TbIII, 21 sextets, 224 quartets and 490 doublets for DyIII, 35 quartets and 112 doublets for ErIII. The spin–orbit coupling based on atomic mean field approximation (AMFI) was taken into account using RASSI-SO with the following numbers of multiplets: 7 septets, 110 quintets, 180 triplets and 180 singlets for TbIII, 21 sextets, 128 quartets and 130 doublets for DyIII, 35 quartets and 112 doublets for ErIII. The relativistic effects were treated with the Douglas–Kroll Hamiltonian. The following basis sets were employed: Ln-ANO-RCC-VQZP (Ln = Tb, Dy and Er for 1–3), O-ANO-RCC-VQZP, N-ANO-RCC-VQZP, C-ANO-RCC-VDZ and H-ANO-RCC-VDZ. Then, the SINGLE_ANISO module was used to calculate all relevant information and magnetic data.

![Fig. 1](https://example.com/fig1.png)

**Fig. 1** (A) Molecular structure of the [Tb(Li)(H2O)(NO3)2]+ complex cation in 1. The thermal ellipsoids are drawn with the 50% probability level. The hydrogen atoms are omitted for clarity. (B) The muffin-like coordination geometry of the [TbN6O4] core in 1.
While many lanthanide complexes show interesting luminescence properties, the luminescence spectra of all the complexes are summarized in Table S2 (ESI†).

Photoluminescence properties

Thus, the coordination number of the LnIII atom in both complexes is nine. The geometries of the coordination polyhedra of lanthanide ions in 1 and 2 are analysed by the program SHAPE 2.1.37 The lowest value of deviation was found for a muffin shape (Table S1†), with the basal trigonal plane formed by O2, N1, and N3 atoms in 1 (O1, N1, and N2 in 2), the equatorial pentagonal plane (O1, O4, N2, N4, and N5 in 1; O2, O4, N3, N4, and N5 in 2), and an O3 atom at the vertex of the muffin (Fig. 1B and S2B†).

The Tb-N bond lengths are in the range of 2.510(2) to 2.671(2) Å and Tb-O bonds vary from 2.347(2) to 2.496(2) Å in 1. The Dy-N bond lengths are in the range of 2.500(3) to 2.659(3) Å, while the Dy-O bonds vary from 2.331(2) to 2.489(2) Å.

Crystal structures of 1 and 2 are stabilized by networks of strong O-H-O hydrogen bonds and weak non-covalent C-H-O and C-H-N interactions. The O-H-O hydrogen bonds connect the coordinated water molecules and non-coordinated nitrate anions, with the O-O separations of 2.746(3) and 2.660(3) Å for 1. Selected non-covalent contacts are summarized in Table S2 (ESI†).

Table 2 Selected interatomic distances (Å) and angles (°) in complexes 1 and 2.

| Distances     | 1               | 2               |
|---------------|-----------------|-----------------|
| Ln-N1         | 2.510(2)        | 2.500(3)        |
| Ln-N2         | 2.599(2)        | 2.567(2)        |
| Ln-N3         | 2.579(2)        | 2.591(3)        |
| Ln-N4         | 2.671(2)        | 2.522(3)        |
| Ln-N5         | 2.528(2)        | 2.639(3)        |
| Ln-O1         | 2.4847(18)      | 2.489(2)        |
| Ln-O2         | 2.4961(18)      | 2.479(2)        |
| Ln-O3         | 2.3468(19)      | 2.331(2)        |
| Ln-O4         | 2.3611(19)      | 2.348(2)        |
| Angles        |                 |                 |
| O3-Ln-N1      | 135.01(7)       | 135.08(9)       |
| O3-Ln-N2      | 77.87(7)        | 146.84(9)       |
| O3-Ln-N3      | 146.96(8)       | 77.74(9)        |
| O3-Ln-N4      | 75.74(7)        | 81.08(9)        |
| O3-Ln-N5      | 81.35(7)        | 75.89(8)        |
| O3-Ln-O1      | 73.90(6)        | 132.74(8)       |
| O3-Ln-O2      | 132.85(6)       | 73.68(8)        |
| O3-Ln-O4      | 81.97(7)        | 81.82(8)        |

Magnetic properties

Static magnetic properties. The temperature and field dependent static magnetic data were acquired on polycrystalline samples of 1–3 as shown in Fig. 2. The value of \( \mu_{\text{eff}}/\mu_{\text{B}} \) at room temperature is 9.68 for 1, 10.39 for 2, and 9.16 for 3, and is close to the expected paramagnetic value of 9.72 (TbIII, \( 7F_6 \)), 10.65 (DyIII, \( 6H_{15/2} \)), and 9.58 (ErIII, \( 4I_{13/2} \)). There is a gradual decrease of \( \mu_{\text{eff}}/\mu_{\text{B}} \) upon cooling the samples to 1.9 K for all compounds 1–3, which is due to depopulation of ligand field multiplets arising from ground atomic terms effected by spin-orbit coupling and a ligand field of the chromophores. There

Photoluminescence properties

While many lanthanide complexes show interesting luminescence properties, the luminescence spectra of all the complexes 1–3 were measured. In accordance with the literature, a reasonable signal was obtained only in the case of complex 1. Its photoluminescence spectrum in the acetonitrile/methanol mixture (\( c = 1 \times 10^{-3} \) mol dm\(^{-3} \)) recorded at room temperature exhibited a broad minor peak at 306 nm assignable to the \( \pi \rightarrow \pi^* \) transition and intense peaks with the maxima at 492, 547, 587 and 622 nm, attributable to the \( ^5D_4 \rightarrow ^7F_6, ^5D_4 \rightarrow ^7F_5, ^5D_4 \rightarrow ^7F_4 \), and \( ^5D_4 \rightarrow ^7F_3 \) transitions, which are typically observed in spectra of TbIII complexes.38 In order to investigate an “antenna effect” of the 2-pyridylmethyl pendant arms in L, a comparison of 1 with the TbIII complex with a parent macrocycle 15-ppyN\(_3\)O\(_2\)\(^{19}\) without pendant arms was performed. As is shown in Fig. S3,† complex 1 exhibits a strong characteristic emission in the visible region. This observation is in accordance with the fact that the pyridine moiety is well known for sensing of lanthanide emission.38,40

Fig. 2 Magnetic data of complexes 1–3. Temperature dependence of the effective magnetic moment (left) and the isothermal molar magnetizations measured at 2 (∙), 5 (□) and 10 (○) K (right). The full lines correspond to ab initio CASSCF calculations done with MOLCAS/ SINGLE_ANISO, and were scaled by the factors: \( f = 0.990 \) for 1, \( f = 0.947 \) for 2 and \( f = 0.877 \) for 3.
are no maxima on susceptibility, which excludes the existence of significant intermolecular contacts of the antiferromagnetic nature. The reciprocal susceptibilities were analysed using the Curie–Weiss law in the temperature range of 25–300 K (Fig. S4†), which resulted in $C = 1.49 \times 10^{-4}$ m$^3$ mol$^{-1}$ K, $\Theta = -4.9$ K and $g = 1.50$ for 1, $C = 1.73 \times 10^{-4}$ m$^3$ mol$^{-1}$ K, $\Theta = -7.9$ K and $g = 1.31$ for 2, $C = 1.38 \times 10^{-4}$ m$^3$ mol$^{-1}$ K, $\Theta = -16.3$ K and $g = 1.17$ for 3. All the Weiss constants are of negative values and $g$-factors are close to theoretical Landé $g$-factors, i.e. 1.50, 1.33, and 1.20 for Tb$^{III}$, Dy$^{III}$, and Er$^{III}$, respectively. The isothermal magnetization data, $M_{\text{mol}}/N_{\text{A}}B_{DC}$, measured at $T = 2$ K saturate to 5.1 for 1, 5.6 for 2 and 4.6 for 3 and these values are well below theoretically predicted values based on $J$ and Landé $g$-factors, which are 9.0 for Tb$^{III}$, 10.0 for Dy$^{III}$ and 9.0 for Er$^{III}$. This points out to large magnetic anisotropy of these complexes.

Dynamic magnetic properties

In order to examine the possible SMM properties of the herein studied coordination compounds 1–3, the ac susceptibility measurements were performed first in zero and nonzero static magnetic fields as depicted in Fig. S5†. None of the compounds showed a nonzero out-of-phase signal of ac susceptibility at zero static magnetic field, but evidently, a small magnetic field must be applied to observe slow relaxation of magnetization and suppression of the tunneling effect. Therefore, the temperature dependence of ac susceptibility was measured at $B_{\text{DC}} = 0.1$ T for frequencies of 1–1500 Hz as shown in Fig. 3. Only in the case of the Dy$^{III}$ compound 2, we observed clearly maxima of out-of-phase susceptibility dependent of the applied frequency and these data were then analysed with the one-component Debye model

$$\chi(\sigma) = \chi_S + (\chi_T - \chi_S)/[1 + (i\sigma\tau)^{1-a}]$$ \hspace{1cm} (1)

which resulted in isothermal ($\chi_T$) and adiabatic ($\chi_S$) susceptibilities, relaxation times ($\tau$) and distribution parameters ($a$) (Table S3†) and construction of the Argand (Cole–Cole) plot (Fig. S6f). Then, the Arrhenius equation was applied to the temperature dependence of the relaxation times, which resulted in the relaxation time $\tau_0 = 2.63 \times 10^{-8}$ s and the effective magnetization reversal barrier $U = 24.4$ K (16.9 cm$^{-1}$) – Fig. S6f.† The ac susceptibility data for 1 and 3 cannot be analysed with eqn (1) due to the absence of maxima on imaginary susceptibility. Therefore, we used a simplified model† according to eqn (2)

$$\ln(\chi'/\chi) = \ln(2\sigma f\tau_0) + U/kT$$ \hspace{1cm} (2)

where higher temperature ac data for higher applied frequencies were included as shown in Fig. S7.† The linear regression analysis resulted in sets of parameters listed in Table 3. The variations in the fitted parameters refer to the distributions of relaxation processes, which are reflected in eqn (1) by parameter $a$. Maximal $U$ were found as $U = 43.5$ K for 1, $U = 64.1$ K for 2 and $U = 82.2$ K for 3. In the case of 2, $U = 64.1$ K is 2.7 times larger than $U = 24.4$ K derived from eqn (1), which can be explained by the fact that the analysis based on eqn (1) is limited only to ac susceptibility data having maxima in the Argand diagram, which means data measured between $T = 1.9$ and 2.8 K. However, a non-zero out-of-phase ac susceptibility is already observed below 4 K (Fig. 3) and especially high temperature data should correspond to the Orbach relaxation mechanism, thus eqn (2) could lead to a better estimate of the relaxation barrier.

Theoretical calculations

In order to better understand the magnetic properties of these compounds, the CASSCF calculations were performed using

![Fig. 3](View Article Online)
MOLCAS 8.0 together with the SINGLE_ANISO program on complexes [Ln(L)(NO₃)(H₂O)](NO₃)₂, 1–3 using geometries following from experimental X-ray data. The resulting wave functions and the energies of the molecular multipoles were used for the calculation of the magnetic properties, g tensors of the lowest Kramers doublet states (Tables S4–S6† Fig. S8†). In the case of Tb³⁺ compound 1, a priori there are no Kramers doublets arising from atomic multiplet T₁₂g, however, the two lowest energy states are almost degenerate, and treating them as Kramers doublets with Sₚ₀ = 1/2 resulted in gₓ = 16.7 and gᵧ = 0.00 (Table S4†), which means that there is large axial magnetic anisotropy. In the case of Kramers doublet ions, Dy³⁺ and Er³⁺, we were able to construct a scheme of the magnetization blocking barrier as shown in Fig. 4. The values displayed on each arrow are the mean absolute values for the corresponding matrix elements of the transition magnetic moment and for values larger than 0.1 an efficient relaxation mechanism is expected. Evidently, the tunnelling mechanism is probable in both compounds 2 and 3, and also thermal relaxation through the first excited state. The coefficients are slightly lower in the case of 2, which probably explains slower relaxation of the magnetization in the case of the Dy³⁺ compound in spite of the larger U in the Er³⁺ compound. The maximal U values extracted from ac susceptibility data with eqn (2), U ≈ 43.5 K for 1, U = 64.1 K for 2 and U = 82.2 K for 3, follow the trend from the calculated energies of the first excited state within this series by CASSCF, 33.3 K for 1, 55.0 K for 2 and 69.0 K for 3 (Tables S4–S6†). Also, the magnetic properties were calculated with the SINGLE_ANISO module, and are compared to the experimental ones in Fig. 2.

Conclusions

In summary, we have successfully prepared three mononuclear lanthanide complexes (Tb³⁺ 1, Dy³⁺ 2, and Er³⁺ 3) with a macrocyclic ligand \{(3,12-bis(2-pyridylmethyl)-3,12,18-triaza-6,9-dioxabicyclo[12.3.1]octadeca-1,14,16-triene; L\} containing two 2-pyridylmethyl pendant arms, and these complexes were characterized structurally and magnetically. In all the cases, a central lanthanide(m) atom was coordinated by a N₅O₄-donor set of L, but due to the high ligand flexibility, the coordination sphere was completed by one water molecule and one nitrate ligand. Thus, the lanthanide(m) atom revealed a coordination number of nine with the N₅O₄-donor set, with a muffin-like geometry. The ac susceptibility measurements showed that all three compounds behave as field-induced single-molecule magnets with the estimated energy barriers U ≈ 44–82 K. The CASSCF calculations analysed with the SINGLE_ANISO module were helpful in understanding this behaviour, and the reasons for that are as follows: (i) the same increasing trend (Tb → Dy → Er) for the first excited state energy gap was found as resulted from the analysis of ac susceptibility; (ii) the lowest transition probabilities for the relaxation of the magnetization were found for 2, for which several maxima out-of-phase ac susceptibilities were already observed; (iii) a relatively small energy splitting and large non-collinearity of the easy axes explain why this series of SMMs shows fast relaxation of magnetization. To summarize, the herein reported compounds 1–3 are the first lanthanide-based SMMs comprising a 2-pyridylmethyl pendant-armed macrocyclic ligand. Nevertheless, the herein reported 15-membered macrocyclic ligand L was proved to be prospective for the synthesis of magnetically interesting coordination compounds, and the investigation of alternative arm-groups is underway.

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