We present the AGILE gamma-ray observations in the energy range 50 MeV–10 GeV of the supernova remnant (SNR) W44, one of the most interesting systems for studying cosmic-ray production. W44 is an intermediate-age SNR (∼20,000 years) and its ejecta expand in a dense medium as shown by a prominent radio shell, nearby molecular clouds, and bright [S II] emitting regions. We extend our gamma-ray analysis to energies substantially lower than previous measurements which could not conclusively establish the nature of the radiation. We find that gamma-ray emission matches remarkably well both the position and shape of the inner SNR shocked plasma. Furthermore, the gamma-ray spectrum shows a prominent peak near 1 GeV with a clear decrement at energies below a few hundreds of MeV as expected from neutral pion decay. Here we demonstrate that (1) hadron-dominated models are consistent with all W44 multwaveshort constraints derived from radio, optical, X-ray, and gamma-ray observations; (2) ad hoc lepton-dominated models fail to explain simultaneously the well-constrained gamma-ray and radio spectra, and require a circumstellar density much larger than the value derived from observations; and (3) the hadron energy spectrum is well described by a power law (with index $\gamma = 3.0 \pm 0.1$) and a low-energy cut-off at $E_c = 6 \pm 1$ GeV. Direct evidence for pion emission is then established in an SNR for the first time.

**Key words:** acceleration of particles – cosmic rays – gamma rays: general – ISM: supernova remnants

**Online-only material:** color figures

1. INTRODUCTION

Providing an unambiguous proof of the cosmic-ray origin until now has been elusive, despite many decades of attempts and controversial claims (e.g., Fermi 1949; Ginzburg & Syrovatskii 1964; Torres et al. 2003; Aharonian 2004; Berezhko & Voelk 2007; Butt 2009). Cosmic rays are mainly protons and heavy ions (hadrons) and, in a few percent, electrons and positrons. Supernova remnants (SNRs) are ideal candidates for the cosmic-ray production up to energies near $E_{\text{kin}} = 10^{15}$ eV. The SNR energy output in the Galaxy can indeed supply the energy budget necessary to maintain the present population of cosmic rays. Furthermore, the observations of ultra-relativistic electrons support the hypothesis that protons are also accelerated in these objects (for a recent review, see Reynolds 2008 and references therein). Proving the fact that the SNR origin of hadronic cosmic rays is difficult because of the complexity of the SNR–environment interaction. From an observational point of view, a direct proof can be given by an unambiguous detection of the gamma-ray emission expected from neutral pion decay in hadronic interactions. However, radiation from spatially accelerated electrons can mask and sometimes over-
to provide an energy spectrum starting at 50 MeV for bright objects (Vercellone et al. 2009; Giuliani et al. 2010a; Vittorini et al. 2011). In this paper, we report on a low-energy γ-ray and multiwavelength spectrum for the SNR W44 in order to constrain the emitting particle spectrum and discriminate between leptonic and hadronic models.

2. THE SUPERNova REMNANT W44

SNR W44 (G34.7−0.4) is a well-studied middle-aged (~20,000 yr) SNR located in the Galactic disk at a distance of ~3 kpc from Earth (Clark & Caswell 1976; Wolszczan et al. 1991). W44 is an ideal system to test the presence of accelerated hadrons and the interplay between hadronic and leptonic models. Radio (Castelli et al. 2007 and references therein) and X-ray (Watson et al. 1983) mapping of the SNR show a roughly elliptical shocked shell and a centrally peaked emission. In the IR band (Reach et al. 2005) it is shown that the shell is expanding into a dense surrounding medium (n ~ 100 cm⁻³). Wolszczan et al. (1991) discovered the radio pulsar PSR B1853+01 with distance and age compatible with the SNR. Wootten (1977) and then Rho et al. (1994) found a massive molecular cloud (MC) interacting with the southeastern side of the remnant. Evidence for a more complex system of massive MCs and for their interactions with the remnant, shown by some features typical of a strong shock, was given by Seta et al. (2004) and then by Reach et al. (2005). The MC–SNR interactions were confirmed by the maser OH (1720 MHz) emission reported by Claussen et al. (1997) and then by Hoffman et al. (2005).

The first estimation of the spectral radio index variations as a function of position over the remnant was done by Castelli et al. (2007): the eastern limb spectrum was consistent with a diffusive shock acceleration model and the spectrum flattening in the westernmost arc confirmed the MC–SNR interaction observed in IR and optical band.

Gamma-ray emission from this SNR has been detected by the Fermi/LAT instrument at energy E > 200 MeV (Abdo et al. 2010a), suggesting the presence of accelerated protons interacting with the surrounding medium.

3. DATA ANALYSIS

AGILE-GRID data were analyzed using the AGILE Standard Analysis Pipeline. We used γ-ray events filtered by means of the F.M.3.119.2 AGILE Filter Pipeline (as described in Vercellone et al. 2008). In order to discriminate between background events and gamma rays, the GRID and anticoincidence system (ACS) signals are processed, reconstructed, and selected by a dedicated software (Giuliani et al. 2006). We used the most recent versions of the diffusion model (Giuliani et al. 2004) and of the calibration files, available at the ASDC site (http://www.asdc.asi.it). We created counts, exposure, and Galactic background gamma-ray maps with a bin size of 0.02 × 0.02. In order to derive the source average flux and spectrum we ran the AGILE point-source analysis software ALIKE (Bulgarelli et al. 2011, based on the maximum likelihood technique described in Mattox et al. 1993) over the whole observing period 2007 July to 2011 April. Both statistic and systematic uncertainties are taken into account. The spectrum was obtained by computing the γ-ray flux in six energy bins selected with the aim to have a significance σ > 4. In order to study the source morphology, we obtained an intensity map integrated over the energy range where the GRID angular resolution is optimal (E > 400 MeV).

4. RESULTS AND DISCUSSION

AGILE detects SNR W44 with a significance of 15.8σ as an extended source. Figure 1(a) shows the AGILE gamma-ray intensity map above 400 MeV of the W44 region with the 324 MHz VLA radio contours. The gamma-ray morphology remarkably resembles the quasi-elliptical pattern of the interior of the radio shell, especially coinciding with the radio brightness enhancements toward the northwest and southeast regions. Both the radio pulsar PSR B1803+01 position (Petre et al. 2002) and its (small) pulsar wind nebula (Giacani et al. 1997, and references therein) are inconsistent with the gamma-ray morphology detected by AGILE. Moreover, Abdo et al. (2010a) excluded the presence of a pulsation in the γ-ray signal. Figure 1(b) shows the CO emission at a kinematic velocity compatible with the distance of W44, tracing the presence of MCs in the W44 surroundings, together with the gamma-ray contour levels. It can be inferred that the southeastern side of the γ-ray source overlaps with the MC–SNR interaction region. In the northern part of the shell, the γ-ray and CO emissions are not correlated, however many studies of the surrounding interstellar medium showed the presence of dense gas not traced by CO (Reach et al. 2005). Figure 1(c) displays an S 14CO map overlaid with X-ray and gamma-ray contours. The strong sulfur [S II] emission, along with Hα emission, indicates the presence of shocked gas (Draine & McKee 1993).

4.1. The Gamma-ray Spectrum

Figure 2 shows the AGILE W44 photon energy spectrum for the whole range 50 MeV–10 GeV. The measured flux above 400 MeV is F = (16.0 ± 1.2) × 10⁻³ photons cm⁻² s⁻¹. In this figure are also shown Fermi/LAT spectral points (Abdo et al. 2010a) over an energy range 0.2–30 GeV. In the band where the spectra overlap, the sets of data are compatible at 1σ. In this regard it is important to stress how AGILE is able to detect the emission from W44 in the energy range 50 MeV–300 MeV extending the spectrum to energies substantially lower than those previously obtained. This spectrum shows a clear decrement at photon energies lower than 400 MeV, confirming the expectations based on neutral pion emission from accelerated protons/ions: a peak energy near 1 GeV for hadron energy spectra flatter than E⁻² (e.g., Aharonian 2004). The gamma-ray spectrum becomes steep at higher energies with a photon power-law index α ~ 3 ± 0.1, in agreement with previous measurements (Abdo et al. 2010a).

4.2. Lepton-dominated Models of Emission

Leptonic-only models of gamma-ray production have to satisfy the very well-determined spatial and spectral constraints provided by the radio, optical, and gamma-ray emissions. For both the Bremsstrahlung and inverse Compton cases we tested the class of leptonic-only models attempting to reproduce the observed gamma-ray spectrum for different input leptonic spectra (see Table 1). We used a phenomenological approach based on two evidences: (1) the synchrotron spectrum implies that the radio electron distribution is well described by a power law over a wide range of energies, (2) the discontinuity in the slope of the gamma-rays spectrum implies a discontinuity in the emitting particle spectrum. We assumed three electron

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17 As pointed out by Rho et al. (1994) in their X-ray and optical study of W44, the optical filaments and the X-ray image showing locally bright emission clumps along the filaments suggest that both are produced by the interaction between the SN shock front and regions of enhanced ambient density.
distributions that can, in principle, describe this behavior: (1) a power law with a high-energy cut-off, $F_1(E) = K_e E^{-\gamma_1} e^{-E/E_c}$ (as in Hendrick & Reynolds 2001), (2) a power law with a low-energy cut-off, $F_2(E) = K_e E^{-\gamma_2} e^{-E_c/E}$ (as in Gabici et al. 2009), (3) a broken power law, $F_3(E) = K_e (E/E_c)^{\gamma_1}(1 + E/E_c)^{\gamma_2-\gamma_1}$ (as in Zirakashvili et al. 2007), where $E_c$ is the cut-off energy and $K_e$ is the normalization constant.

Since the distributions with a cut-off (1 and 2) fail to reproduce simultaneously both the radio and $\gamma$-ray spectrum, we refer to the distribution (3). The best fit to the gamma-ray data is obtained with $E_c = 1$ GeV, and indices $\gamma_1 = 0$ and $\gamma_2 = 3.3$ above and below $E_c$, respectively. Synchrotron emission originating from this distribution can be evaluated for different values of the average magnetic field. Figure 3 shows the case of the most “favorable” leptonic-only model characterized by $B = 20$ $\mu$G (other cases turn out to be even less favorable). We find that, at high frequencies, the calculated synchrotron spectrum is in strong disagreement (factor larger than 4) with the radio emission produced co-spatially to the gamma-ray one (Castelletti et al. 2007, Figure 1). Furthermore, the inferred average density, $n = 300$ cm$^{-3}$, is too large (by a factor of three) compared with the circumstellar medium constraints (Reach et al. 2005); a lower value for $n$ would be incompatible with the gamma-ray spectrum fitting. Similar or even stronger contradictions with the multwavelength data apply to other leptonic-only models that we systematically explored for a large variety of parameters.
In the case of inverse Compton dominated models, two sources of soft photons are available: the cosmic background radiation (CBR) and the interstellar radiation field (ISRF). In the first case, a second peak in the gamma-ray spectrum is unavoidably expected with a peak energy $E_{\text{max}} \sim 1$ TeV, in contradiction with the upper limits obtained from TeV Cherenkov telescopes. In the case of interaction with ISRF, instead, the calculated synchrotron peak is not compatible with the radio continuum data for any reasonable value of the magnetic field in the SNR shell. We can then reliably exclude leptonic-only models of emission for SNR W44.

4.3. Hadron-dominated Models of Emission

It is interesting to determine the main physical parameters of an emission model dominated by hadrons in the gamma-ray energy range. The gamma-ray emission was derived assuming that protons interact with the nuclei of the ambient medium through $pp$ interactions and then radiate through $\pi^0$ decay (Kelner et al. 2006). In order to fit the gamma-ray spectrum we tested, as proton energy distribution, a power law, a broken power law, and a power law with cut-offs. The distribution providing the best-fit of the gamma-ray spectrum turned out to be a power law with a relatively steep spectral index and a low-energy cut-off, $F_p(E) \sim E^{-p} e^{-E_p/E}$, with $p = 3.0 \pm 0.1$ and $E_p = 6 \pm 1$ GeV. The inferred total energy of hadrons producing the observed gamma-ray spectrum is $E_h = 10^{49}$–$10^{50}$ erg (depending on the average magnetic field and local densities, see Table 1) corresponding to a fraction of the total SNR shell kinetic energy $\epsilon \sim 0.01$–0.1. This value is a lower limit on the total energy of hadrons accelerated by W44 during its entire life, because the most energetic hadrons, likely accelerated during the early epochs of the SNR life, can escape from the acceleration site on timescales shorter than the age of this SNR (Berezinskii et al. 1990; Gabici et al. 2009). Electrons produce the radio continuum and shell-like features by synchrotron emission in an average magnetic field in the range $B \sim 10$–$100 \mu G$. They also produce Bremsstrahlung and inverse Compton components that are subdominant in the gamma-ray energy range for a standard value of the electron/proton number ratio $\chi \sim 0.01$. Our best “hadronic model” is shown in Figure 4. We note that the spectral index found in this analysis is quite steep in comparison with the expectations for the spectrum of particles accelerated in SNRs (e.g., Reynolds 2008) and that the cosmic-ray spectrum can be altered (softened) by the interaction with sites where

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**Table 1**

| Model          | $(n)$ (cm$^{-3}$) | $(B)$ ($\mu G$) | $E_h$ (erg) | $E_c$ (erg) | $p_1$ | $p_2$ | Comments                                      |
|----------------|------------------|-----------------|-------------|-------------|-------|-------|-----------------------------------------------|
| Hadronic       | 20               | 20              | $2.1 \times 10^{20}$ | $2.0 \times 10^{49}$ | 1.7   | 3.0   | Good agreement with data                      |
|                | 100              | 70              | $3.3 \times 10^{29}$ | $2.8 \times 10^{48}$ | 1.7   | 3.0   | Good agreement with all data                  |
| Leptonic       | 300              | 20              | ...          | $7.3 \times 10^{18}$ | 0     | 3.3   | Incompatible with the spectrum                |
|                | 5000             | 200             | ...          | $4.4 \times 10^{17}$ | 0     | 3.3   | Incompatible with the radio spectrum: wrong slope, low-frequency excess |

(A color version of this figure is available in the online journal.)

In the case of inverse Compton dominated models, two sources of soft photons are available: the cosmic background radiation (CBR) and the interstellar radiation field (ISRF). In the first case, a second peak in the gamma-ray spectrum is unavoidable with a peak energy $E_{\text{max}} \sim 1$ TeV, in contradiction with the upper limits obtained from TeV Cherenkov telescopes. In the case of interaction with ISRF, instead, the calculated synchrotron peak is not compatible with the radio continuum data for any reasonable value of the magnetic field in the SNR shell. We can then reliably exclude leptonic-only models of emission for SNR W44.

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Figure 4. Theoretical modeling of the broadband spectrum of SNR W44 superimposed with the radio (data points in red color) and gamma-ray data of Figure 2 (in blue color). TeV upper limits are also shown. This is a hadronic scenario characterized by $B = 70 \mu G$ and $n = 100 \, \text{cm}^{-3}$. The yellow curve shows the neutral pion emission from the accelerated proton distribution discussed in the text. The green curves show the electron contribution by synchrotron (dashed curve), bremsstrahlung (solid curve), and IC (dotted curve) emissions. The red curve shows the total gamma-ray emission.

(A color version of this figure is available in the online journal.)

the target gas is confined, for example, by MCs (see, e.g., Gabici et al. 2010; Ohira et al. 2011). A comprehensive theoretical interpretation of this fact is beyond the scope of this work and will be addressed in a forthcoming paper.

5. CONCLUSIONS

W44 turns out to be an extremely interesting SNR whose environment and shocked material configuration favor a detailed testing of hadronic versus leptonic emission. We find that models of accelerated protons/ions interacting with nearby dense gas successfully explain all the observed multifrequency properties of W44. The total number of accelerated hadrons in W44 is consistent with an efficiency of a few percent in the conversion of SNR kinetic energy into cosmic-ray energy.

Within the hadronic scenario, the broadband spectrum is successfully modeled by the gamma-ray emission resulting from the decay of neutral pions produced by accelerated protons/ions with a spectral index near 3 and a cut-off energy $E_\text{c} \sim 6 \text{ GeV}$ (Figures 2 and 4). Electrons (a few percent in number) contribute to the synchrotron radio emission and to a weak bremsstrahlung component. It is interesting to note that the spectral cut-off energy $E_\text{c}$ and the steep power-law index may indicate either a “diffusion” effect of high-energy hadrons, or a suppression of efficient particle acceleration in dense environments (Uchiyama et al. 2010; Malkov et al. 2011) as expected in W44. Our results on W44 are complementary with those obtained in other SNRs, most notably IC 443 (Tavani et al. 2010) and W28 (Giuliani et al. 2010b). Also in these SNRs the complex interaction of cosmic rays with their dense surroundings produces gamma-ray spectra with high-energy cut-offs observed in the range 1–10 GeV.

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