THE $^7$Be II RESONANCE LINES IN TWO CLASSICAL NOVAE V5668 SGR AND V2944 OPH

AKİTO TAJITSU$^1$, KOZO SADAKANE$^2$, HIROYUKI NAITO$^3$, AKIRA ARAI$^4$, HIDEYO KAWAKITA$^4$, AND WAKO AKI$^5$

$^1$ Subaru Telescope, National Astronomical Observatory of Japan, 650 North A‘ohoku Place, Hilo, HI 96720, USA; tajitsu@naoj.org
$^2$ Astronomical Institute, Osaka Kyoiku University, Asahigaoka, Kashiwara, Osaka 582-8582, Japan
$^3$ Yayiro Observatory, 157-1 Nishin, Yayiro, Hokkaido 066-0066, Japan
$^4$ Koyama Astronomical Observatory, Kyoto Sangyo University, Motoyama, Kamigamo, Kita-ku, Kyoto 603-8555, Japan
$^5$ National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan

Received 2015 October 15; accepted 2016 January 13; published 2016 February 18

ABSTRACT

We report spectroscopic observations of the resonance lines of singly ionized $^7$Be in the blueshifted absorption line systems found in the post-outburst spectra of two classical novae—V5668 Sgr (Nova Sagittarii 2015 No. 2) and V2944 Oph (Nova Ophiuchi 2015). The unstable isotope $^7$Be should have been created during the thermonuclear runaway (TNR) of these novae and decayed to form $^7$Li within a short period (a half-life of 53.22 days). These confirmations of $^7$Be are the second and the third ones following the first case found in V339 Del by Tajitsu et al. The blueshifted absorption line systems in both novae are clearly divided into two velocity components, both of which contain $^7$Be. This means that the absorbing gases in both velocity components consist of products of TNR. We estimated the amounts of $^7$Be produced during the outbursts of both novae and concluded that significant $^7$Li should have been created. These findings strongly suggest that the explosive production of $^7$Li via the reaction $^4$He($\alpha$, $\gamma$)$^7$Be and its subsequent decay to $^7$Li occurs frequently among classical novae and contributes to the process of Galactic Li enrichment.

Key words: Galaxy: evolution – novae, cataclysmic variables – nuclear reactions, nucleosynthesis, abundances – stars: individual (V5668 Sgr, V2944 Oph)

1. INTRODUCTION

Lithium (Li) is a key element in the study of the chemical evolution of the universe because it has likely been produced in various sites and events—the Big Bang nucleosynthesis, interactions between energetic cosmic rays and interstellar matter, evolved low-mass stars, novae, and supernova explosions. The observed Li evolutionary curve has a plateau for young Galactic ages (<2.5 Gyr) followed by a steep rise. This indicates that relatively low-mass stellar components with long lifetimes are major sources of Li in the recent universe (Romano et al. 1999, 2001; Prantzos 2012). Some low-mass evolved stars have indeed been found to have Li-enriched surfaces (e.g., Melo et al. 2005). However, there has been no observational confirmation that such stars supply Li to interstellar medium. Li contained in stellar surfaces could easily be depleted by convection because Li should be destroyed inside stars where the temperature is higher than 2.5 million K.

Novae, which evolve from low-mass binaries, are expected to be one of the candidates for Li suppliers in the recent universe because they experience explosive nucleosynthesis and mass-loss simultaneously. Cameron & Fowler (1971) made the first theoretical study of Li production in novae. They predicted that the radioactive isotopes of beryllium, $^7$Be, was produced via the reaction $^4$He($\alpha$, $\gamma$)$^7$Be during the thermonuclear runaway (TNR) in the accumulated thin gas layer on the surface of a white dwarf (WD). The $^7$Be that is produced will be blown away from the surface of a WD by the outburst wind and then decay to form $^7$Li in the cooler interstellar material within a short period of time (a half-life of 53.22 days). In the 1990s–70s, details of this process, which is called the Cameron-Fowler process, were studied by theoretical analyses (Arnould & Nörgaard 1975; Starrfield et al. 1978; Boffin et al. 1993; Hernandez et al. 1996; José & Hernanz 1998).

Recently we have reported the detection of $^7$Be in the post-outburst UV spectra of the classical nova V339 Del (Tajitsu et al. 2015). Because the resonance doublet lines of Be II are located at ~3130 Å where the telluric absorption that is mainly caused by ozone severely obstructs observations from ground-based telescopes, quantitative studies of stellar Be abundances were enabled in 1990s by the advent of ~10 m class telescopes on high mountains. Such telescopes are equipped with high-resolution spectrographs with high sensitivities in the near-UV region (e.g., Boesgaard et al. 1999 (Keck I telescope); Primas et al. 2000 (VLT)). Extensive analyses of stellar Be abundances were conducted in large samples of metal-poor stars (Boesgaard et al. 2011) and solar analogs (Takeda et al. 2011). In the case of V339 Del we found that the absorption lines of Be II were purely originating from $^7$Be instead of the stable $^9$Be. Therefore it was the first observational confirmation of the Cameron-Fowler process in classical novae. The $^7$Be was found in highly blueshifted (radial velocities, $v_{\text{rad}} \sim 1000$ km s$^{-1}$) absorption lines, which are assumed to correspond to nova ejecta blown off the surface of a WD by the outburst wind. Through the direct comparison of absorption strengths between $^7$Be and Ca II we estimated that the $^7$Be abundance in the ejecta of V339 Del could be higher than the pre-existing theoretical prediction (José & Hernanz 1998). Izzo et al. (2015) reported the presence of the Li I λ6708 line in the blueshifted absorption line system in the early phase (7–13 days after the outburst) spectra of V1369 Cen (Nova Centauri 2013). The overabundance of Li in this nova amounts to the order of 10$^4$ with respect to the solar photospheric composition.

Now it is quite interesting to know how commonly this $^7$Be ($=^7$Li) production occurs among various types of classical novae and to evaluate their contribution to the Li evolution in the Galaxy. Prompted by these interests we have performed high-resolution spectroscopic observations of two Galactic...
The Astrophysical Journal, 818:191 (9pp), 2016 February 20

TAJITSU ET AL.

novae discovered in March 2015—V5668 Sgr and V2944 Oph—during their early declining phases.

V5668 Sgr (Nova Sagittarii 2015 No. 2 = PNV J18365700−2855420) was discovered as a bright 6.0 mag (unfiltered) source by John Seach on 2015 March 15.634 UT and announced in the American Association of Variable Star Observers (AAVSO) Alert Notice (Waagen 2015a). The nova recorded its optical maximum ($V = 4.3$) on March 21.67 UT (MJD = 57102.67). Its optical magnitudes stayed close to the maximum ($V \sim 4.5–6.5$) for about 80 days then showed a rapid decline by dust formation (Banerjee et al. 2015a, 2015b; Gehrz et al. 2015). The possible candidates of the progenitor were found in the USNO-B1.0 catalog (Monet et al. 2003) within $6''$ from the nova (USNO-B1.0 0610−078494.3, $B1 = 17.17$, $R1 = 16.82$, $R2 = 16.42$, and $I = 15.55$ mag; USNO-B1.0 0610−0784923, $R1 = 14.36$, $B2 = 13.62$, $R2 = 13.65$, and $I = 15.08$ mag) and in the Galaxy Evolution Explorer (GALEX; Martin et al. 2005) catalog of NUV (1171–2831 Å) sources 3′′6 from the nova (GALEX J183656.8−285539, NUV = 17.323 mag) (Waagen 2015a). From a spectroscopic observation obtained immediately after the discovery, Williams et al. (2015) reported that it is an Fe ii type nova according to the classification introduced by Williams (1992).

V2944 Oph (Nova Ophiuchi 2015 = PNV J17291350−1846120) was discovered as a 12.2 mag (unfiltered) source by Yukio Sakurai on 2015 March 29.766 UT (Waagen 2015b). After its discovery the nova showed a gradual rise to its optical maximum at $V = 9.0$ on April 14.75 UT (MJD = 57126.75). Then it showed a quick decline to $V \sim 11.5$ within the following few days and stayed at $V \sim 12$ for >80 days. In the USNO-B1.0 catalog a possible candidate for the progenitor was found as a ∼18 mag star at the position of the nova (Waagen 2015b). Its spectrum during the pre-maximum phase showed He/N type characteristics (Danilet et al. 2015). However, Munari et al. (2015) concluded that it is an Fe ii type nova by spectroscopic observations carried out during the maximum and early decline phase. Munari & Walter (2016) presented observations of the time evolution in line profiles during the pre-maximum phase and pointed out the presence of pre-existing circumstellar materials around the nova. They concluded that the companion in the nova system is not a main-sequence dwarf as in most novae but rather an evolved subgiant.

In this paper we report the detection of strong $^7$Be absorption lines in these two novae. In Section 2 we describe our high-resolution spectroscopic observations and data analysis. The results from our observations are described in Section 3. The $^7$Be abundances estimated from our observed data are presented in Section 4. A brief discussion and conclusions are given in Section 5.

2. OBSERVATIONS

The post-outburst spectra of the two classical novae were obtained using the High Dispersion Spectrograph (HDS) (Noguchi et al. 2002) of the 8.2 m Subaru Telescope on 2015 May 29 (V5668 Sgr; +69 days after the optical maximum) and July 3 (V2944 Oph; +80 days), respectively. Figure 1 shows the AAVSO light curves of two novae with the epochs of our HDS observations. In the case of V5668 Sgr our observation was just before the start of the rapid decline in optical magnitudes by dust formation. For each nova we obtained spectra under three configurations of the spectrograph, which cover the wavelength regions from 3030 to 4630 Å, from 4110 to 6860 Å, and from 6670 to 9360 Å. Spectral resolving power was set to $R \approx 60,000$ and 45,000 with 0.′′6 (0.3 mm) and 0.′′8 (0.4 mm) slit widths, respectively. The details of observations are reported in Table 1. Data reduction was carried out using the IRAF software in a standard manner. The nonlinearity in the detectors are corrected by the method described in Tajitsu et al. (2010). The typical residual of wavelength calibration performed using Th–Ar comparison spectrum is smaller than 10$^{-3}$ Å for each spectrograph configuration. Spectrophotometric calibration was performed using the spectrum of σ Sgr ($V = 2.058$, B2V) and was obtained nearly at the same altitude of the nova on the same nights. All spectra were converted to the heliocentric scale. Correction for interstellar extinction has not been applied. The line identifications were carried out using the atomic line database of Kramida et al. (2014). We also used the line list of Kurucz & Bell (1995) for some weak transitions of Fe-peak elements.
3. RESULTS

The post-outburst spectra of V5668 Sgr (+69 days) and V2944 Oph (+80 days) displayed in Figure 2 and Figure 3, respectively, show very similar spectral characteristics. They contain a series of strong broad emission lines originating from HI Balmer series and some other permitted transitions originating from Fe II. As in the case of V339 Del described in Tajitsu et al. (2015), most of these broad emission lines are accompanied by blueshifted multiple absorption lines whose radial velocities correspond to ∼−2000 to −800 km s⁻¹. Many blueshifted absorption lines of Ti II, Cr II, and Fe II are found in the near-UV region.

Table 1

| Date        | UT  | MJD     | Exposure | Range    | Resolution |
|-------------|-----|---------|----------|----------|------------|
| V5668 Sgr   |     |         |          |          |            |
| May 29      | 10  | 57,171.447 | 600     | 4110–6860 | 60,000     |
| (+69 days)b | 12  | 57,171.502 | 3600    | 3030–4630 | 60,000     |
|             | 13  | 57,171.581 | 300     | 6670–9360 | 60,000     |
| V2944 Oph   |     |         |          |          |            |
| Jul 03      | 6   | 57,206.289 | 900     | 6670–9360 | 45,000     |
| (+80 days)c | 7   | 57,206.299 | 900     | 4110–6860 | 45,000     |
|             | 7   | 57,206.322 | 5400    | 3030–4630 | 45,000     |

Notes.

a UT is the universal time at the start of an exposure.
b Days after the optical (V) maximum (MJD = 57,102.67).
c Days after the optical (V) maximum (MJD = 57,126.75).

Figure 2. Example sections of the flux-calibrated HDS spectrum of V5668 Sgr obtained at +69 days. Enlarged views of the spectrum in the hatched areas are presented in Figure 4. Identified absorption line systems originating from Fe-peak elements are indicated by ticks at the top.
Figure 4 displays the enlarged views of the spectrum of V5668 Sgr in the vicinities of the H\(\gamma\), Ca\(\text{II}\) K, \(^7\)Be\(\text{II}\) \(\lambda\)3131.228, Na\(\text{I}\) D2, Fe\(\text{II}\) (the multiplet number 42; Moore 1959) \(\lambda\)5169.03, and Li\(\text{I}\) lines on the velocity scale. The blueshifted absorption lines of the H\(\gamma\) line is clearly found to be divided into two components: the low-velocity component (LVC) is a single sharp absorption line at a relatively low \(v_{\text{rad}}\approx -786\) km s\(^{-1}\) and the high-velocity component (HVC) contains a series of subcomponents spreading between \(-2200 < v_{\text{rad}} < -1350\) km s\(^{-1}\). The radial velocities of each subcomponent are measured at the absorption dips in the Fe\(\text{II}\) (42) and H\(\gamma\) line and indicated with vertical lines in the figure. The sharp absorption of the LVC can be easily identified in all lines in the figure. These LVC lines originating from singly ionized Fe-peak elements (Fe\(\text{II}\), Ti\(\text{II}\), Cr\(\text{II}\), Mn\(\text{II}\), and Ni\(\text{II}\)) dominate the spectrum in the UV range as in the case of V339 Del from \(+38\) to \(+48\) days. All of the identified transitions originate from levels of low excitation potentials (the lower energy level of the transition, \(E_{\text{lower}} \lesssim 3.1\) eV). The absorption strengths of the HVC vary among different transitions. All of six subcomponents in the HVC can be identified in the H\(\gamma\) and Fe\(\text{II}\) (42) lines, although some of them are weaker or missing in Ca\(\text{II}\) and Na\(\text{I}\) lines. In the UV range only the subcomponents at \(v_{\text{rad}} = -1385\) and \(-1444\) km s\(^{-1}\) of the HVC can be identified in the strong lines originating from Fe-peak elements.

Figure 5 presents the spectrum of V2944 Oph displayed in the same way as in Figure 4. The spectrum of this nova shows quite similar characteristics to that of V5668 Sgr. It also shows the blueshifted absorption systems divided into two velocity components: the LVC at \(v_{\text{rad}} = -878\) km s\(^{-1}\) and the HVC between \(-2000 < v_{\text{rad}} < -1300\) km s\(^{-1}\). At least six subcomponents can be identified in the HVC as shown by vertical lines in the figure.

As shown in the bottom panels of Figures 4 and 5, we cannot find any counterparts of blueshifted absorption line systems of the \(^7\)Li\(\text{I}\) \(\lambda\)6708 in spite of the high signal-to-noise ratios (S/N) of the spectra (\(~700\) and \(~160\) at 6690 Å in V5668 Sgr and V2944 Oph, respectively). There is also no counterpart of the Ca\(\text{I}\) (32) \(\lambda\)6718 in the spectra of both novae, although this line has been detected together with the Li\(\text{I}\) \(\lambda\)6708 line in the first three weeks spectra of V1369 Cen (Izzo et al. 2015). The Ca\(\text{I}\) \(\lambda\)4227 and K\(\text{I}\) \(\lambda\)7699 lines do not have any counterparts in the LVC and the HVC in both observed novae. These observations imply that the ejected gas had been heated into a state so hot that almost all Li and Ca atoms are ionized in the circumstance during the epochs of our HDS observations.

Figure 6 displays the enlarged radial velocity profiles of Fe\(\text{II}\), H\(\gamma\), and the \(^7\)Be\(\text{II}\) “doublet” lines in both novae. Although the \(^7\)Be\(\text{II}\) \(\lambda\lambda\) 3130.583 and 3131.228 doublet lines coalesce, the dips of each line can be partially identified. In V5668 Sgr, two absorption dips (A and B) are clearly found in all lines. In
V2944 Oph, two dips (B and C) of the $^7$Be II doublet coincide with those of $H\gamma$. The dip A of $^7$Be II $\lambda$3131.228 (the weaker component of the doublet) in the LVC in this nova is unclear due to the poor S/N, which is partly disturbed by cosmic ray.

Figure 4. Blueshifted absorption line systems in the spectrum of V5668 Sgr. The spectrum obtained at +69 days is displayed in the vicinities of $Hy$, Ca II K, $^7$Be II $\lambda$3131.228, Na i D2, Fe II $\lambda$5169.03, and $^7$Li i $\lambda$6707.76 lines (top to bottom) on the velocity scale (upper horizontal). The local continua, fitted with high-order (15–25) spline functions, are overplotted with thick lines (orange in the online version). Dips of individual absorption line, which can be identified in Fe II and/or $H\gamma$, are indicated with vertical lines (green in the online version). The identified dips by the other component of the doublet are also indicated in the $^7$Be II and Na i D panels. In the bottom panel no counterparts of the blueshifted absorption line systems of the $^7$Li i or the Ca i (32) lines are found in their expected wavelengths (solid and dashed vertical lines).

Figure 5. Same as Figure 4 but for the spectrum of V2944 Oph obtained at +80 days. The spectrum is smoothed by a three-pixel boxcar. Cosmic ray hits on the spectrum in the vicinity of $^7$Be II are indicated with crosses. In the bottom panel there are no counterparts of the $^7$Li i doublet or the Ca i (32) lines.

V2944 Oph, two dips (B and C) of the $^7$Be II doublet coincide with those of $H\gamma$. The dip A of $^7$Be II $\lambda$3131.228 (the weaker component of the doublet) in the LVC in this nova is unclear due to the poor S/N, which is partly disturbed by cosmic ray.
hits on the spectrum. A subcomponent in the LVC at \( v_{\text{rad}} = -849 \) km s\(^{-1}\) (the dip A+) is found in the Fe\(\Pi\) and both \(^7\text{Be}\Pi\) lines. We can find no alternative candidates of identification for these dips from line lists of Fe-peak elements. Because no dips corresponding to the \(^9\text{Be}\Pi\) doublet [the isotopic shift; \( \Delta \lambda = -0.161 \) Å (Yan et al. 2008)] can be found, we conclude that the absorption systems at \( \sim 3110 – 3125 \) Å in both novae purely originate from \(^7\text{Be}\Pi\) as in the case of V339 Del. The dips of the \(^7\text{Be}\Pi\) doublet become hard to be identified toward the flat bottoms of the HVC in both novae. This indicates that the absorption of \(^7\text{Be}\Pi\) is saturated while the intensities at the flat bottoms remain \( \sim 25 \) and \( \sim 15\% \) of the continua in V5668 Sgr and V2944 Oph, respectively. In V5668 Sgr the HVCs of the Balmer lines show wavy bottom shapes. This difference indicates that the \(^7\text{Be}\Pi\) absorption is stronger than the Balmer lines at this velocity range.

It is interesting to see the similarity in the velocity profiles of the blueshifted absorption systems in both novae. Some classical novae also have similar blueshifted absorption systems with two distinct velocity components—the narrow LVC and the broad HVC (e.g., Nova LMC 2005, V2574 Oph in Williams et al. 2008), which are known as the “principal” and the “diffuse enhanced” systems in post-maximum nova spectra (McLaughlin 1960; Munari 2014). V1369 Cen at +13 days also has two velocity components at \( v_{\text{rad}} = -550 \) and \( -1300 < v_{\text{rad}} < -900 \) (Izzo et al. 2015). One remarkable point found in this study is that \(^7\text{Be}\) is detected in both the LVC and the HVC. This means that both of the absorption systems consist of nova ejecta which have experienced TNR. The characteristics of the LVCs in V5668 Sgr and V2944 Oph closely coincide with those of the Transient Heavy Element Absorption (THEA) found by Williams et al. (2008) in post-outburst spectra of classical novae. They

---

Figure 6. Enlarged radial velocity profiles of the LVC and the red side edge of the HVC in V5668 Sgr (upper) and V2944 Oph (lower). Normalized spectra in the vicinities of Fe\(\Pi\)(42)\(\lambda 5169\), H\(\gamma\), and the \(^7\text{Be}\Pi\) doublet are displayed with offsets. The solid vertical lines indicate absorption components identified in the \(^7\text{Be}\Pi\) doublet and other lines. The dashed vertical lines show absorption components identified only in Fe\(\Pi\) and/or Balmer lines. The expected positions of the \(^9\text{Be}\Pi\) doublet are indicated by partial dashed lines. In the LVC in V2944 Oph (lower bottom panel) the spectrum in the vicinity of \(^7\text{Be}\Pi\) lines is smoothed by a three-pixel boxcar, and a subcomponent at \( v_{\text{rad}} = -849 \) km s\(^{-1}\) is found in each line (the dip A+; the orange vertical line in the online version).
assumed that these THEAs are produced by pre-outburst ejecta coming from the secondary star. Considering our results on the $^7$Be detection, however, it is natural to assume that the THEAs are produced by nova ejecta that have experienced TNR.

The $^7$Be II absorption may be contaminated by LVCs originating from other Fe-peak elements. We set requirements for major contaminants as $E_{\text{lower}} < 3.1$ eV and $\log (gf) > −1.0$ and select candidates for contaminants as listed in Table 2. Only some Cr II (5) lines are found to be the main contaminants to the LVC and the HVC of the $^7$Be II doublet. In V5668 Sgr, the dip of the Cr II (5)$\lambda 3132.053$ is clearly identified within the $^7$Be II LVC (Figure 6 upper right panel). In the $^7$Be II HVC, no Cr II (5) lines can be identified by the strong saturation effect. However, their equivalent widths ($W$) can be estimated using measured equivalent widths of nearby unblended Cr II lines belonging to the same multiplets. We estimate that the combined contributions of these contaminants are less than 5% of the total equivalent widths of the $^7$Be II absorption in both the LVC and the HVC in V5668 Sgr. In the case of V2944 Oph we find that only the Cr II (5)$\lambda 3132.053$ is present besides the LVC of the $^7$Be II doublet. We estimate that the combined contribution of contaminants in the HVC in this nova is similar to or less than that in V5668 Sgr.

4. $^7$Be ABUNDANCE

We placed “best effort” local continua along each line ($−3000 < \nu_{\text{obs}} < 1000$ km s$^{-1}$) on the complex, undulating spectra of two novae by the fitting with a high-order (15–25) spline function to measure the equivalent widths of the absorption lines. Although the $^7$Be II doublet cannot be resolved, its total equivalent widths, $W(^7\text{Be} \text{II})$, in the LVC and the HVC are measured. After subtracting the combined effects of contaminations discussed above we tabulate the resulting $W(^7\text{Be} \text{II})$ in Table 3. The errors of $W$ are estimated from the S/N of the neighboring continuum. They do not include the uncertainties in the continuum placements.

The abundances of $^7$Be in nova ejecta is estimated following the same procedure as used in Tajitsu et al. (2015). We compare the equivalent widths of Ca II K ($\log (gf) = +0.135$) and $^7$Be II $\lambda 3130.583 + \lambda 3131.228$ ($\log (gf) = −0.178, −0.479$, respectively) in the LVC, which show a simpler structure than the HVC. Here we need the following assumptions: (1) both the $^7$Be II and Ca II lines are not saturated and (2) the covering factor of the absorbing gas cloud to the background illuminating source has no wavelength dependence. We notice that the LVC feature of $^7$Be II at $−786$ km s$^{-1}$ of V5668 Sgr might be saturated. The bottom of the $^7$Be II lines in the LVC almost reaches to that in the HVC where its flat profile suggests that the absorption is saturated. The corresponding LVC component of the Ca II K line in V5668 Sgr is shallow and the line appears to be unsaturated. The abundance of $^7$Be in V5668 Sgr obtained under such a condition should be taken as a lower limit.

Following Spitzer (1998, Equations (3)-(48)), the ratio of column number densities, $N$, can be written as

$$
\frac{N(^7\text{Be} \text{II})}{N(\text{Ca II})} = \frac{W(^7\text{Be} \text{II}) \lambda(\text{Ca II} K)^2}{\lambda(^7\text{Be} \text{II})^2 W(\text{Ca II} K)} \times \frac{gf(\text{Ca II} K)}{gf(^7\text{Be} \text{II} \lambda 3130.583) + gf(^7\text{Be} \text{II} \lambda 3131.228)}
$$

$$
= \frac{W(^7\text{Be} \text{II})}{313^2} \times \frac{3934^2}{W(\text{Ca II} K)} \times \frac{10^{+0.135}}{10^{-0.178} + 10^{-0.479}}.
$$

Using data of the $W(^7\text{Be} \text{II})$ and the $W(\text{Ca II} K)$ listed in Table 3 the column density ratios $N(^7\text{Be} \text{II})/N(\text{Ca II})$ derived in the LVC of V5668 Sgr and V2944 Oph are 8.1 ± 2.0 and 2.5 ± 1.5, respectively. These ratios are close to that found in V339 Del at $+47$ days.

The ionization states of each species are keys to convert the above ratios into atomic abundances. We assume that the most of $^7$Be is in the singly ionized state because it still stays in hot ejecta. We also assume that Ca is in the singly ionized state. The absence of Ca I lines in the blueshifted absorption line systems supports this assumption. The difference in the second ionization potential (Be: 18.21 eV, Ca: 11.87 eV) may lead to different balances between the second and the third ionized ions and this may be a source of an error in determining the abundances. We tried but failed to detect absorption lines of doubly ionized ions of Fe-peak elements, which have intermediate second ionization potentials, consulting Kramida et al. (2014). Absorption lines of Fe II are observed in optical

---

**Table 2**

| Lines       | $\lambda_{\text{obs}}$ | $\log (gf)$ | $\lambda_{\text{LVC}}^a$ | $^7$Be II $^b$ | $W$ (Å)$^c$ | $\lambda_{\text{LVC}}^a$ | $^7$Be II $^b$ | $W$ (Å)$^c$ |
|-------------|------------------------|-------------|--------------------------|----------------|-------------|--------------------------|----------------|-------------|
| Ti II (67)  | 3106.246               | −0.170      | 3098.11                  | ...            | <0.003      | 3097.16                  | ...            | <0.006      |
| Ti II (67)  | 3117.676               | −0.500      | (3109.51) HVC            | ...            | (3109.51) HVC | (3109.52) HVC            | (0.037 ± 0.020) | ...         |
| Cr II (5)   | 3118.646               | 0.000       | (3110.48) HVC            | (0.054 ± 0.026)| (3110.65) HVC | (3111.23) HVC            | (0.049 ± 0.030)| ...         |
| Cr II (5)   | 3119.825               | −0.460      | (3111.65) HVC            | ...            | (3110.69) HVC | (3111.36) HVC            | (0.036 ± 0.020)| ...         |
| Cr II (5)   | 3120.359               | +0.120      | (3112.18) HVC            | (0.071 ± 0.034)| (3112.32) HVC | (3111.23) HVC            | (0.049 ± 0.030)| ...         |
| Cr II (5)   | 3120.497               | −0.018      | (3112.32) HVC            | (0.052 ± 0.025)| (3120.49) ... | (3120.49) ...            | ...            | ...         |
| Cr II (5)   | 3128.692               | −0.320      | ...                      | 0.013 ± 0.003  | ...          | 0.013 ± 0.003            | ...            | ...         |
| Cr II (5)   | 3132.053               | +0.079      | 3123.85                  | LVC            | 0.067 ± 0.015| 3123.85                  | ...            | 0.045 ± 0.015|
| Cr II (5)   | 3136.681               | −0.250      | 3128.46                  | ...            | 0.045 ± 0.015| 3128.46                  | ...            | <0.006      |

**Notes.**

$^a$ Observed and expected wavelengths of the LVCs ($\nu_{\text{obs}} = −786$ and $−878$ km s$^{-1}$ in V5668 Sgr and V2944 Oph, respectively) of each transition. The undetected lines in the $^7$Be II HVC are listed in brackets.

$^b$ Velocity component of the $^7$Be II doublet expected to be contaminated by the line.

$^c$ Observed or expected equivalent widths of each line. The values estimated using nearby same multiplets are listed in brackets.
range spectra of B-type stars (Thompson et al. 2008). This may support the assumption that all of $^7$Be and Ca contained in the LVC are in the singly ionized state. Adopting this assumption, the mass fraction of $^7$Be to the sum of all constituent mass components, $X(\text{Be} \ \|)$, can be presented as: $X(\text{Be} \ \|) \sim N(\text{Be} \ \|)/N(\text{Ca} \ \|) \times 7/40 \times X(\text{Ca})$. We have obtained $X(\text{Be} \ / \ X(\text{Ca}) = 1.4 \pm 0.3$ for V5668 Sgr and $0.44 \pm 0.26$ for V2944 Oph. These $X(\text{Be} \ / \ X(\text{Ca})$ values correspond to logarithmic overabundances of Li by $+6.2$ dex and $+5.7$ dex with respect to the solar photospheric abundance (Asplund et al. 2009). As in the case of V339 Del, the Li abundances in ejecta of both nova explosions are comparable to or even higher than those of Ca. Because the results presented here are based on the observations $\sim$75–100 days after their outbursts, the amount of $^7$Be freshly produced in the TNR could be three to four times larger than our measurements.

We note that these abundances are obtained only using a limited fraction of the nova ejecta giving rise to the LVC absorption. They do not necessarily represent the abundances in the whole materials erupted from these novae. Furthermore, when the $^7$Be II lines are saturated the above $X(\text{Be} \ / \ X(\text{Ca})$ values should be treated carefully. Therefore, the derived $^7$Be abundances in this study might involve large uncertainties. However, the above result implies that classical novae are indeed playing an important role in the process of the Galactic Li enrichment.

To produce $^7$Be via the reaction $^3$He($\alpha$, $\gamma$)$^7$Be the total number of $^3$He atoms in the accreted gas before TNR should be larger than or equal to that of freshly produced $^7$Be atoms in the nova ejecta. This means that the number ratios, $N(\text{He})/N(\text{Ca})$, in the accreted gases of V5668 Sgr and V2944 Oph are $\geq 8.1 \pm 2.0$ and $\geq 25.5 \pm 1.5$, respectively. If we adopt the solar abundance ratio between He and Ca (log (N(He)/N(Ca)) = $+4.59$; Asplund et al. 2009), the lower limits of the number ratios between $^3$He and $^7$Be, $N(\text{He})/N(\text{He})$, in the accreted gases of V5668 Sgr and V2944 Oph should amount to 0.020 $\pm 0.005$ and 0.006 $\pm 0.003\%$, respectively. These isotopic ratios are close to that found in the solar system ($=0.0166\%$; Asplund et al. 2009). Hernanz et al. (1996) pointed out that $^3$He destruction via the reaction $^3$He($\alpha$, $\gamma$)$^7$He occurs concurrently with $^3$He($\alpha$, $\gamma$)$^7$Be during TNR. They noted that the $^3$He destruction via $^3$He($\alpha$, $\gamma$)$^7$He is less pronounced in CO nova than in ONe nova. Because both V5668 Sgr and V2944 Oph are found to be CO novae (see Section 5) we guess that the process of $^3$He destruction should be ineffective and that most of the $^3$He in the accreted gases has been converted into $^7$Be in both novae.

5. DISCUSSION AND CONCLUSIONS

The detection of $^7$Be in the post-outburst spectra of V5668 Sgr and V2944 Oph strongly suggests that the explosive production of $^7$Be is a popular phenomenon among classical novae. The light curves of both novae show slower evolution than that of V339 Del in which the rapid decline was observed in $\sim$+40 days. Adopting the relation between the timescaling factor of the light curve and the WD mass (e.g., Hachisu & Kato 2006, 2015), the WD masses of both novae are lower than that of V339 Del. It is natural to assume that both novae have CO WDs as in the case of V339 Del. Therefore we infer that the explosive Li production via the Cameron-Fowler process occurs frequently at least among CO novae.

It is interesting to notice that we have found LVC and HVC absorptions of the Na I D lines in V5668 Sgr and V2944 Oph while no trace of blueshifted absorptions of the Na I D lines were detected in V339 Del (Tajitsu et al. 2015). We find no absorption features corresponding to Li I $\lambda$6708, K I $\lambda$7699, and Ca I $\lambda$4227 in both V5668 Sgr and V2944 Oph as well as in V339 Del. Izzo et al. (2015) found blueshifted absorptions of H$\beta$ and Na I D lines in V1369 Cen at two velocities ($\sim$1100 and $\sim$550 km s$^{-1}$) on spectra obtained during an early phase after the outburst ($7$–$13$ days). They identified Li I $\lambda$6708, K I $\lambda$7699, Ca I $\lambda$6718, and Ca I $\lambda$4227 in the lower-velocity component. These apparent spectral differences such as detection or no detection of absorptions of the Na I D or Ca I lines should reflect differences in physical conditions (most probably the ionization state) in the absorbing clouds. On the other hand, the presence of $^7$Be both in the LVC and the HVC strongly indicates that the absorbing clouds in both velocity components consist of materials that have experienced TNR.

The abundance of Li obtained by Izzo et al. (2015) from the Li I $\lambda$6708 is lower than the $^7$Be (=Li) abundances obtained in the present study. There had been no report in the literature of detection of the Li I $\lambda$6708 line among post-outburst novae before Izzo et al. (2015). Because the Li I $\lambda$6708 is located in the easily accessible red spectral region, the paucity of observational information may imply that ejecta of post-outburst novae rarely permit the physical condition in which neutral Li can survive. However, it will be interesting to carry out long-term monitoring observations of post-outburst spectra of novae including both $^7$Be and Li lines and examine the consistency of derived Li abundances. Further observations of the $^7$Be II line in various types of classical novae will help to quantify roles of novae as the source of Li in the current universe.

This work is based on data observed at the Subaru Telescope, which is operated by the National Astronomical Observatory of Japan (NAOJ). We acknowledge with thanks the variable star observations from the AAVSO International Database contributed by observers worldwide and used in this research.

Facility: Subaru(HDS).
REFERENCES

Arnould, M., & Nörgaard, H. 1975, A&A, 42, 55
Asplund, M., Grevesse, N., Sauval, A. J., & Scott, P. 2009, ARA&A, 47, 481
Banerjee, D. P. K., Ashok, N. M., & Srivastava, M. 2015a, ATel, 7299, 1
Banerjee, D. P. K., Ashok, N. M., Venkataraman, V., & Srivastava, M. 2015b, ATel, 7303, 1
Boesgaard, A. M., Deliyannis, C. P., King, J. R., et al. 1999, AJ, 117, 1549
Boesgaard, A. M., Rich, J. A., Levesque, E. M., & Bowler, B. P. 2011, ApJ, 743, 140
Boffin, H. M. J., Paulus, G., Arnould, M., & Mowlavi, N. 1993, A&A, 279, 173
Cameron, A. G. W., & Fowler, W. A. 1971, ApJ, 164, 111
Danilet, A. B., Holoien, T. W.-S., Wagner, R. M., et al. 2015, ATel, 7339, 1
Gehrz, R. D., Evans, A., Woodward, C. E., et al. 2015, ATel, 7862, 1
Hachisu, I., & Kato, M. 2006, ApJS, 167, 59
Hachisu, I., & Kato, M. 2015, ApJ, 798, 76
Hernanz, M., Jose, J., Coc, A., & Isern, J. 1996, ApJL, 465, L27
Izzo, L., Della Valle, M., Mason, E., et al. 2015, ApJL, 808, L14
Jose, J., & Hernanz, M. 1998, ApJ, 494, 680
Kramida, A.,Ralchenko, Y., Reader, J. & NIST ASD Team 2014, NIST Atomic Spectra Database (ver. 5.2) (Gaithersburg, MD: National Institute of Standards and Technology) available:http://physics.nist.gov/asd
Kurucz, R., & Bell, B. 1995, Atomic Line Data, Kurucz CD-ROM No. 23 (Cambridge, MA: Smithsonian Astrophysical Observatory)
Martin, D. C., Fanson, J., Schiminovich, D., et al. 2005, ApJL, 619, L1
McLaughlin, D. B. 1960, in Stellar Atmospheres, ed. J. L. Greenstein (Chicago, IL: Univ. Chicago Press), 585
Melo, C. H. F., de Laverny, P., Santos, N. C., et al. 2005, A&A, 439, 227
Monet, D. G., Levine, S. E., Canzian, B., et al. 2003, AJ, 125, 984
Moore, C. E. 1959, A Multiplet Table of Astrophysical Interest: NBS Technical Note No. 36 (Washington, D.C.: US Department of Commerce) Reprinted Version of the 1945 Edition
Munari, U. 2014, in ASP Conf. Ser. 490, Stella Nova: Past and Future Decades, ed. P. A. Woudt, & V. A. R. M. Ribeiro (San Francisco, CA: ASP), 183
Munari, U., Ochner, P., Siviero, A., Dallaporta, S., & Valisa, P. 2015, ATel, 7367, 1
Munari, U., & Walter, F. M. 2016, MNRAS, 455, L57
Noguchi, K., Aoki, W., Kawanomoto, S., et al. 2002, PASJ, 54, 855
Prantzos, N. 2012, A&A, 542, A67
Primas, F., Molaro, P., Bonifacio, P., & Hill, V. 2000, A&A, 362, 666
Romano, D., Matteucci, F., Molaro, P., & Bonifacio, P. 1999, A&A, 352, 117
Romano, D., Matteucci, F., Ventura, P., & D’Antona, F. 2001, A&A, 374, 646
Spitzer, L. 1998, Physical Processes in the Interstellar Medium (Weinheim: Wiley-VCH)
Starrfield, S., Truran, J. W., Sparks, W. M., & Arnould, M. 1978, ApJ, 222, 600
Tajitsu, A., Aoki, W., Kawanomoto, S., & Narita, N. 2010, PNAOJ, 13, 1
Tajitsu, A., Sadakane, K., Naito, H., Arai, A., & Aoki, W. 2015, Natur, 518, 381
Takeda, Y., Tajitsu, A., Honda, S., et al. 2011, PASJ, 63, 697
Thompson, H. M. A., Keenan, F. P., Dufour, P. L., et al. 2008, MNRAS, 383, 729
Waagen, E. 0. 2015a, AAN, 512, 1
Waagen, E. 0. 2015b, AAN, 516, 1
Williams, R., Mason, E., Della Valle, M., & Ederoclite, A. 2008, ApJ, 685, 451
Williams, R. E. 1992, AJ, 104, 725
Williams, S. C., Darnley, M. J., & Bode, M. F. 2015, ATel, 7230, 1
Yan, Z.-C., Nörtershäuser, W., & Drake, G. W. F. 2008, PhRvL, 100, 243002