Evidence for Shock-heated Gas in the Taffy Galaxies and Bridge from Optical Emission-line IFU Spectroscopy

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Received 2018 December 18; revised 2019 April 28; accepted 2019 May 11; published 2019 June 25

Abstract

We present optical integral field unit observations of the Taffy system (UGC 12914/15), named for the radio emission that stretches between the two galaxies. Given that these gas-rich galaxies are believed to have recently collided head-on, the pair exhibits a surprisingly normal total (sub-LIRG) IR luminosity ($L_{\text{IR}} \sim 4.5 \times 10^{10} L_{\odot}$). Previous observations have demonstrated that a large quantity of molecular and neutral gas has been drawn out of the galaxies into a massive multiphase bridge. We present, for the first time, spatially resolved spectroscopy of the ionized gas in the system. The results show that the ionized gas is highly disturbed kinematically, with gas spread in two main filaments between the two galaxies. The line profiles exhibit widespread double components in both the bridge and parts of the disks of the galaxies. We investigate the spatial distribution of the excitation properties of the ionized gas using emission-line diagnostic diagrams and conclude that a large quantity (up to 40%) of the emission from the entire system is consistent with gas heated in $\sim 200 \text{ km} \text{s}^{-1}$ shocks. While the shocked gas is mainly associated with the bridge, there is a significant amount of shocked gas associated with both galaxies. Confirming other multiwavelength indicators, the results suggest that the effects of shocks and turbulence can continue to be felt in a high-speed galaxy collision long after the collision has occurred. The persistence of shocks in the Taffy system may explain the relatively low current star formation rates in the system as a whole.

Key words: galaxies: individual (UGC 12914/5) – galaxies: interactions – galaxies: ISM – galaxies: kinematics and dynamics

1. Introduction

It is now generally accepted that collisions and mergers between gas-rich galaxies often generate intense star formation activity and associated strong infrared emission (e.g., Joseph & Wright 1985; Soifer et al. 1987; Soifer & Neugebauer 1991). Ultraluminous infrared galaxies (ULIRGs; $L_{\text{IR}} > 10^{12} L_{\odot}$) and luminous infrared galaxies (LIRGs; $10^{11} > L_{\text{IR}}/L_{\odot} > 10^{10}$) frequently involve mergers or interactions of gas-rich galaxy pairs, with the likelihood of them being associated with a major merger increasing with infrared (IR) luminosity (Armus et al. 1987, 2009; Sanders et al. 1988a, 1988b; Sanders & Mirabel 1996; Elbaz et al. 2002). While it is clear that major mergers play an important role in generating high IR luminosities in the local universe, their role at higher redshift is still being explored.

Shocks and turbulence potentially play a role in changing the conditions of the gas in collisional galaxies, not always leading to enhancements in star formation. In the local Hickson Compact Groups, Alatalo et al. (2015) found evidence that multiple collisions can quench or significantly suppress star formation in some systems where turbulence and shocks are present (see also Lisenfeld et al. 2017). These galaxies had previously been found to contain large volumes of warm molecular hydrogen that emit their energy mainly in the mid-IR, and were believed to be shock heated (Cluver et al. 2013). An extreme example is found in the Stephan’s Quintet system. Here, a large filament of molecular gas is found in the intergalactic medium, in which a large fraction of the gas is warm and in a shock-heated phase (Appleton et al. 2006, 2017; Guillard et al. 2009; Cluver et al. 2010). Shocks, though hard to detect in LIRGs and ULIRGs because of the dominant effects of star formation on optical emission-line diagnostics, are being increasingly detected with the advent of spatially resolved optical integral field unit (IFU) spectroscopy (Rich et al. 2011, 2014, 2015). How large-scale shocks and turbulence affect the star formation in such galaxies, as well as how important this process is in higher-redshift systems, is currently unknown.

An interesting example of an ongoing major merger that may be caught in a highly disturbed state is the Taffy galaxy pair UGC 12915/4 (hereafter Taffy-N and Taffy-S for simplicity). Despite having recently undergone a strong head-on collision, the Taffy system appears surprisingly normal in its IR properties, with a total $L_{\text{IR}} = 4.5 \times 10^{10} L_{\odot}$ summed over the whole system based on multiwavelength Spitzer and Herschel spectral energy distribution (SED) photometric fitting (Appleton et al. 2015; see also Jarrett et al. 1999; Sanders et al. 2003). The reason that the system is so normal in the IR,
despite its recent violent history, is not known. It may be that we are catching the Taffy system in a peculiar moment where most of its gas is so disturbed that it cannot yet generate significant star formation. If so, studying the conditions of the gas in between the galaxies (referred to as the bridge) may well yield interesting insight into how shocks and turbulence can inhibit star formation in violently colliding galaxies. The Taffy galaxies were named for the discovery of a bridge of radio continuum emission, stretching, like saltwater taffy (candy), between the galaxies (Condon et al. 1993). Evidence suggests that the two galaxies collided 25–30 Myr ago, allowing their stellar components to pass through each other, but stripping \( \sim 7 \times 10^9 M_\odot \) of molecular and atomic gas into a bridge between them (Braine et al. 2003; Gao et al. 2003; Zhu et al. 2007). There is more gas in the bridge than in the two galaxies combined.

The bridge appears to be strongly disturbed (and probably turbulent), based on kinematically broad CO line studies of the bridge and strong mid-IR \( H_2 \) emission and [C II]157.7 \( \mu m \) lines suggestive of shocks (Peterson et al. 2012, 2018). Despite its high gas mass, the average star formation rate (SFR) in the entire bridge through SED fitting is quite low, \( \sim 0.45 M_\odot \text{yr}^{-1} \), excluding the prominent extragalactic H II region seen south-west of UGC 12915, which was separately found to have an SFR of \( 0.24 M_\odot \text{yr}^{-1} \) (Appleton et al. 2015). Numerical models of such a head-on collision between two gas-rich galaxies (e.g., Struck 1997) and a detailed model of the Taffy system (Vollmer et al. 2012) provide strong support for the idea that the gas left behind in the center-of-mass frame of the collision would be highly turbulent and that some would be strongly shock heated. Appleton et al. (2015) detected faint extended soft X-ray emission and several compact point X-ray sources in the bridge, the former being consistent with shock-heated gas that has not had time to completely cool since the collision occurred. Finally, Lisenfeld & Völk (2010) concluded that the radio emission in the bridge could be explained in terms of cosmic rays accelerated in magnetic fields compressed in shocks.

Although the Taffy galaxies have been studied quite extensively at longer wavelengths (Condon et al. 1993; Jarrett et al. 1999), very little work has been done at visible or near-IR wavelengths. Bushouse (1987) presented early digital video camera observations that showed \( H\alpha \) emission from the inner disks of both galaxies and emission from the extragalactic H II regions in the bridge. Pa\( \alpha \) observations from the ground were also made by Komugi et al. (2012). The galaxies show strong disturbances in their optical structure, including rings and loops, and the possible recent onset of star formation in the bridge, including at least one prominent extragalactic H II region and fainter clusters—some of which are seen in archival NICMOS observations from the Hubble Space Telescope (HST; P. N. Appleton et al. 2019, in preparation).

This paper represents the first major study of the ionized gas phase in the Taffy system and bridge. We provide, for the first time, a detailed exploration of both the kinematics and excitation properties of the optical emission-line gas in the Taffy system. The paper is organized as follows: Section 2 describes the observations and methods used in the paper. We describe the fitting process used on the double-line profiles and the gas kinematics through H\( \alpha \) channel maps and velocity field moment maps in Sections 3.1 and 4, respectively. The effects of dust extinction on the measured line fluxes are discussed in Section 3.2. We describe the results from our line diagnostic diagrams in Section 5. In Section 6, we discuss our results on the properties of the ionized gas and its excitation mechanisms through the use of emission-line ratio diagnostic diagrams and comparison with shock models. In Sections 6 and 6.1 we discuss our estimates of the ionized gas fraction from star formation and the SFR in the system. We discuss the evidence for a post-starburst population in the underlying starlight in Section 6.3. In Section 7 we present our conclusions. We assume a comoving distance to the galaxies of 62 Mpc based on a mean heliocentric velocity for the system of 4350 km s\(^{-1}\) and a Hubble constant of 70 km s\(^{-1}\) Mpc\(^{-1}\).

2. Observations, Data Reduction, and Analysis Methods

The IFU data presented in this work were obtained with the VIRUS-P Spectrograph at McDonald Observatory (Hill et al. 2008; Blanc et al. 2010). VIRUS-P is the Visible Integral-field Replicable Unit Spectrograph prototype (now called the George and Cynthia Mitchell Spectrograph, GCMS) mounted on the 2.7 m Harlan J. Smith telescope. The IFU has 246 fibers (each fiber has an angular diameter of \( 4^\prime \) 16 on the sky) with a one-third filling factor. We used several cycles of a three-point dither pattern to completely cover the 2\( ^\prime 8 \) field of view. We used the gratings VP2 and VP4 for our blue and red channel spectra, respectively. VP2 and VP4 have a spectral resolution of \( 1.6 \) and \( 1.5 \) \( \AA \) and covered a range of \( 4700–5350 \) \( \AA \) and \( 6200–6850 \) \( \AA \), respectively. This spectral resolution corresponds to a velocity resolution of \( \sim 100 \) and \( \sim 70 \) km s\(^{-1}\) at the wavelengths of H\( \beta \) in VP2 and H\( \alpha \) in VP4, respectively.

We made observations of the Taffy galaxies on 2012 January 31 and February 2 (blue spectrometer) and 2012 February 1 (red spectrometer) with a total exposure time per dither position of 2200 s (blue) and 1200 s (red). Conditions were photometrically good at the time of the observations with moderate seeing of 1.8–2.5 (less than the diameter of a fiber).

These data were processed using the VACCINE pipeline, which identified and traced each fiber on the CCD chip and performed bias, flat-field, and wavelength calibration (based on lamp spectra) on a fiber-by-fiber basis for the science frames. VACCINE is a Fortran-based reduction package developed for the HETDEX Pilot Survey (Adams et al. 2011) and the VENGA project (Blanc et al. 2010). Cosmic-ray removal was then performed using the IDL routine LA-Cosmic (van-Dokkum 2001).

2.1. Flux Calibration and Cube Building

The flux calibration and pointing refinements and final cube construction of these data were performed in three steps following the methods described by Blanc et al. (2013): (i) Relative spectrophotometric calibration was performed, which is applied to all the fibers. Observations were made in a six-point dither pattern (including several fibers) of the standard star Hz 15 (HIP 21776). An algorithm was used to solve for the position of the star on the fibers and determine the spectrophotometric transformation from native (ADU) units across the spectrum to units of erg s\(^{-1}\) cm\(^{-2}\) \( \AA \)\(^{-1}\). The results were applied to all the fibers irrespective of throughput. This step resulted in a relative flux uncertainty across the band of \( \sim 8\% \). (ii) Astrometry and absolute flux calibration, using a bootstrapping method, were used to effectively cross-correlate a reconstructed image of the galaxy derived by integrating the
light from each fiber, with a calibrated images of Taffy from the Sloan Digital Sky Survey (SDSS; York et al. 2000) in the g and r bands, suitably convolved to the resolution of the VIRUS-P fiber system. This helped refine the astrometry and the assembly of the final cube from the individual observations of each field. The cross-correlation also allowed the spectrum in each fiber to be absolutely scaled to the SDSS band in question. The details of this procedure are given in Blanc et al. (2013). Tests performed in that paper show that the absolute spectrophotometric flux calibration has a typical accuracy of 15%–30%, after taking into account the uncertainties in SDSS calibration and the VIRUS-P relative spectrophotometric accuracy. (iii) A final flux-calibrated 3D spectral cube was created by combining all the various observational pointing frames into a single interpolated cube with resulting $2 \times 2^n^2$ spaxels ($\sim 0.3$ kpc arcsec$^{-1}$ based on the assumed distance of 62 Mpc). These processes were repeated for the red and blue channels, creating final flux-calibrated blue and red spectral cubes.

2.2. Spectral Mapping, and Continuum and Emission-line Fitting

The processing and extraction of astrophysical information from the data cubes was done using a combination of IRAF/PyRAF, IDL, and Python routines. Before beginning our analysis, we smoothed the data cubes spatially, but not in the spectral direction, using a Gaussian kernel with a standard deviation of 1.47 pixels. This was done to boost the signal-to-noise ratio in areas that we were interested in, particularly the Taffy bridge region, which has relatively low signal-to-noise ratio compared to the galaxies. This spatial smoothing effectively reduces the noise in spectra from individual spaxels by a factor of $\sim 2$. We used these spatially smoothed cubes for all of the analysis done in this work.

For fitting each individual spaxel in the IFU data we used the IDL software toolkit LZIFU (LaZy-IFU; Ho et al. 2016). LZIFU automates fitting multiple emission lines superimposed on a continuum for multiple spaxels in each channel and provides 2D maps of continuum and line fluxes, velocities, and velocity dispersion. It is capable of fitting emission lines that are superimposed on deep absorption features and also emission lines with multiple velocity components. Figures 1 and 2 show our fitting results for two individual spaxels that show these features in their spectra. Figure 1 shows a spaxel that has strong H$\beta$ absorption and H$\beta$ emission superimposed on the absorption trough. This spaxel lies very close to the center of Taffy-N. This absorption must be accounted for with the continuum fitting to get accurate emission-line fluxes, as well as an accurate line profile. The emission lines also show evidence for double-line profiles in many positions across the system. As an example, Figure 2 shows a spaxel that lies close to the edge of the extragalactic H$\Pi$ region, which clearly shows two separate velocity components.

LZIFU works by first fitting the continuum with a custom implementation of the PPXF code (Capellari & Emsellem 2004) within LZIFU and then fitting the emission lines after subtracting the continuum model. The continuum fitting uses a set of model templates that are fit to a combined blue+red channel spectrum. LZIFU also accounts for systematic errors in the models and nonstellar contributions to the continuum data by fitting a multiplicative polynomial simultaneously with the continuum models. The emission lines, along with residuals from sky line subtraction, are masked during the continuum fitting process. We also specified the following lines to be fit (and masked during continuum fitting)—H$\beta$, the [O III] $\lambda\lambda$4959, 5007 doublet, [O I] $\lambda$6300, [O I] $\lambda$6364, H$\alpha$, the [N II] $\lambda\lambda$6548, 6583 doublet, and the [S II] $\lambda\lambda$6716, 6731 doublet. The extinction-corrected emission-line fluxes for all the lines detected in different regions of the Taffy system (as defined in Figure 3) are tabulated in Table 1.

We ran LZIFU on the entire IFU data cube for the Taffy galaxies, which is $58 \times 58$ and $59 \times 59$ spaxels with 2227 and 2350 spectral elements for the blue and red channels, respectively. The process of fitting the full cube was not straightforward for several reasons relating to the peculiar kinematics of the Taffy system. The LZIFU software was designed to work best with a galaxy showing slowly varying velocity centroids relatively close to the initial guess for the velocity of the system. In the Taffy system, the velocity range of the emission lines over the whole

![Figure 1](image1.png)

**Figure 1.** Our LZIFU fitting results for a spaxel that has deep H$\beta$ absorption and H$\beta$ emission superimposed on the absorption trough. The top left and right panels show the blue and red data, respectively, along with their model fits. The bottom panels show the corresponding residuals from the fitting. The gray line shows the raw data from the spaxel, and the blue and red lines show the model fits to the respective channels. Note that the figure does not show the full wavelength coverage of the data but instead is focused on showing the relevant absorption and emission features, i.e., H$\beta$ and the [O III] $\lambda\lambda$4959, 5007 doublet in the blue channel and H$\alpha$ and its neighboring [N II] $\lambda\lambda$6548, 6583 doublet lines in the red channel.

![Figure 2](image2.png)

**Figure 2.** Same as Figure 1, but now showing the LZIFU fitting results for another spaxel that displays distinct velocity components more clearly.
system was large, with the emission lines sometimes exhibiting complex behavior, in addition to occasionally being observed superimposed on deep Balmer absorption lines. Furthermore, in a large number of spaxels, we found multiple velocity components that did not always move together as a function of position. As a result, a single set of initial guesses for the various starting parameters did not work for the whole cube, but had to be adjusted spatially to achieve good fits, especially in regions showing double-line profiles. After an iterative process of fitting with different starting parameters tuned to particular regions, we were able to get a consistent set of smoothly varying results across the whole system. We provide the most relevant LZIFU parameters and their starting guesses in Table 2.

3. Results

3.1. Emission-line Gas and H\(\beta\) Absorption within the System

Previous observations of the Taffy system in the visible wavelength range had reported the detection of ionized gas within the galaxies and the extragalactic H II region (Bushouse 1986, 1987). Because of the sensitivity of the VIRUS-P instrument to very faint diffuse emission, we report the presence of ionized gas throughout the Taffy system—both within the galaxies, the so-called extragalactic H II region, and also in the bridge between the galaxies. We detect emission from many lines, including H\(\alpha\), H\(\beta\), the [O III] \(\lambda\lambda 5007, 5007\) doublet, the [N II] \(\lambda\lambda 6548, 6583\) doublet, and the [S II] \(\lambda\lambda 6716, 6731\) doublet lines (see Figures 1 and 2). We also detect strong emission from the atomic oxygen line [O I] \(\lambda 6300\) and sometimes the weaker [O I] \(\lambda 6364\) line in the galaxies and the extragalactic H II region in the bridge. Although emission lines dominate in many of the locations across the system, in some cases H\(\beta\) emission is observed superimposed on a broad absorption trough indicative of a post-starburst population (as seen in Figure 1).

In Figure 3 we present some extracted spectra in several places in the system to provide an overview of the complexity of the kinematics in this system. The spectra show expanded
views of the \([\text{[O III]} \lambda 5007]\) and \(\text{H}\alpha\) lines along the major axis of both galaxies and a sampling of the bridge. As with the previous \(\text{H I}\) investigations of Taffy by Condon et al. (1993) and Braine et al. (2003), the ionized gas spectra along the major axis of Taffy-N (N1–N5) show clear rotation from low to high velocities as one proceeds northward, whereas in Taffy-S (S1–S5) the rotation is also obvious, but in the opposite sense. This confirms the suggestion that the galaxies were counterrotating when they collided. Many of the line profiles are complex and contain multiple components. Of special note are the broad lines in the bridge (especially B1 and B2), as well as complex multicomponent structures in the northwest of both galaxies (N4, N5, S4, and S5). The nucleus of Taffy-S (S2) also shows very broad \([\text{[O III]}]\) emission and weaker broad wings in \(\text{H}\alpha\) (especially when a correction is made for \(\text{H}\alpha\) absorption; see Section 4.2). We also show polygons marking additional regions of interest on the SDSS image in Figure 3. These polygons show regions that are investigated in the emission-line diagnostic diagrams in Section 5. In the emission-line diagnostic diagrams, we investigate the excitation mechanisms for the western regions of Taffy-N and the bridge separately because these regions exhibit peculiar kinematics distinct from the rest of Taffy-N and the bridge (see the detailed discussion in Sections 4.1 and 4.2).

Figure 4(a) shows total \(\text{H}\alpha\) emission contours overlaid on an SDSS \(i\)-band image. The two galaxies and the area near the extragalactic \(\text{H}\II\) region can clearly be distinguished as regions with the brightest \(\text{H}\alpha\) emission, with extended emission spread between the galaxies. For Taffy-S, the brightest \(\text{H}\alpha\) lies on either side of the nucleus with fainter emission from the direction of the nucleus. It is noticeable that there is no obvious ionized gas associated with the faint southern part of the outer ring in Taffy-S, although there may be some associated with the northern part of the ring between the galaxies. Taffy-N has a very different distribution of ionized gas, with a strong concentration in the inner disk and fainter emission extending along the northwest major axis, where it appears to join with bridge material.

Figure 4(b) shows a map of the equivalent width of the \(\text{H}\beta\) absorption across the system. In contrast to the lack of emission lines in the southern ring of Taffy-S, strong stellar absorption is seen there, along the extreme northwesterly edge of the galaxy, and in an extended region northeast of the faint stellar ring. In Taffy-N the strongest absorption is seen at the southeastern end of the major axis in the region outside the main body of \(\text{H}\alpha\) emission, although it extends at lower equivalent width as far as the center of the galaxy. We will discuss the implications of this \(\text{H}\beta\) absorption in Section 6.3.

### Table 1

| Region | \(A_V\) | \(\text{H}\beta\) | \([\text{[O III]} \lambda 5007]\) | \([\text{[O I]} \lambda 6300]\) | \(\text{H}\alpha\) | \([\text{[N II]} \lambda 6583]\) | \([\text{S II]} \lambda 6716]\) | \([\text{S II]} \lambda 6731]\) |
|--------|--------|-------------|-----------------|-----------------|-----------|-----------------|-----------------|-----------------|
| Taffy-N| 2.29   | 3.2 ± 0.23  | 1.98 ± 0.2      | 0.59 ± 0.16    | 9.12 ± 0.16| 4.34 ± 0.14    | 2.24 ± 0.14    | 1.35 ± 0.16    |
| Taffy-S| 0.94   | 1.67 ± 0.13 | 0.67 ± 0.11     | 0.39 ± 0.11    | 4.78 ± 0.11| 2.70 ± 0.1     | 1.38 ± 0.1     | 1.18 ± 0.12    |
| Bridge | 1.89   | 1.42 ± 0.19 | 1.48 ± 0.16     | 0.55 ± 0.14    | 4.05 ± 0.13| 1.89 ± 0.11    | 1.25 ± 0.11    | 0.81 ± 0.13    |

**Note.**

*Values quoted here are from the combined Taffy-N and bridge regions, i.e., including the western parts defined in Figure 3.

#### 3.2. Dust Extinction from the Balmer Decrement

We estimate the extinction caused by dust by examining the line ratio of the Balmer lines, \(\text{H}\alpha\) and \(\text{H}\beta\), referred to as the Balmer decrement, and assume case B recombination. The color excess \(E(B-V)\) is given by

\[
E(B - V) = 1.97 \log_{10} \left( \frac{\text{H}\alpha / \text{H}\beta_{\text{abs}}}{2.86} \right).
\]

The extinction at wavelength \(\lambda\) is related to the color excess by

\[
A_\lambda = k(\lambda) E(B - V).
\]

We assume a reddening curve \(k(\lambda)\) of the form given by Calzetti et al. (2000). Adopting the same method for all the observed lines allows us to correct, spaxel by spaxel, the observed line fluxes for dust extinction to arrive at intrinsic line fluxes.

A map of the dust extinction at visual wavelengths, \(A_V\), derived from the Balmer decrement is shown in Figure 5, and superimposed on the map are CO (1−0) contours from Gao et al. (2003) using data from the Berkeley-Illinois-Maryland-Association interferometer. It can be seen that the dust-extinction map follows reasonably well the CO (1−0) surface density map, with a peak at the center of Taffy-N (as expected since the H\(_2\) surface density is very high there). A high value of extinction is also seen at the southern tip of Taffy-S. This extends beyond the CO column density contours, but we note that the VLA maps of H\(_{\text{I}}\) in this region by Condon et al. (1993) show a peak column density there of \(1.8 \times 10^{21} \text{ cm}^{-2}\), which would imply an \(A_V\) of \(\sim 1\) mag integrated over an \(18 \times 18''\) beam (Güver & Özel 2009).

Interestingly, we observe a low value of \(A_V\) at the center of the Taffy-S and along the northern edge of the spiral arm of that galaxy. The latter result is consistent with a low H\(_{2}\) column there, although the low \(A_V\) in the nucleus may imply different conditions in the gas excitation (deviations from the assumed case B recombination—perhaps due to a low-luminosity active galactic nucleus [LAGN]) or a significantly reduced dust-to-gas ratio there.

Given that dust opacity measurements depend strongly on galaxy inclination (see, e.g., Driver et al. 2007; Unterborn & Ryden 2008), we caution that our dust-extinction estimates for both galaxies should be treated as lower limits and that actual \(A_V\) values could be much higher. For example, Gao et al. (2003) estimate \(A_V\) values could be higher than 10 mag for both galaxies. Our low \(A_V\) values for the Taffy galaxies could be because measuring the dust extinction from the \(\text{H}\alpha\) and \(\text{H}\beta\) lines involved in the Balmer decrement effectively only probes the effects of dust superficially (essentially a “skin” effect; e.g.,
Table 2
LZIFU Parameters Supplied to the Configuration File

| Parameter Name | Value       | Description                                                                 |
|----------------|-------------|------------------------------------------------------------------------------|
| only_1side     | 0           | 0: two-sided data; 1: one-sided data                                         |
| z              | 0.0145      | Redshift                                                                     |
| fit_ran        | [4700, 6855]| Fitting range                                                                |

Continuum Fitting with PPIXF

| Parameter Name | Value       | Description                                                                 |
|----------------|-------------|------------------------------------------------------------------------------|
| mask_width     | 12          | Full width to mask around emission lines defined in lzifu_linelist.pro       |
| cont_vel_sig_guess | [0., 50.]  | Starting guess of deIV and vel_dispersion of continuum (in km s\(^{-1}\)) |
| cont_ebv_guess | 0.1         | Starting guess of ebv                                                        |

Emission Fitting with MPFIT

| Parameter Name | Value       | Description                                                                 |
|----------------|-------------|------------------------------------------------------------------------------|
| fit_dlambda    | 22.         | Full width around line centers to be fitted. (A)                             |
| ncomp          | 2           | Number of component                                                         |
| line_sig_guess | 70          | Initial guess of velocity dispersion for emission line (km s\(^{-1}\))      |
| vdisp_ran      | [−50, 500]  | Velocity dispersion constraints in km s\(^{-1}\)                             |
| vel_ran        | [−600., +600.0] | Velocity constraints in km s\(^{-1}\); 0 is systemic velocity from set.z |

Variation in Initial Guess\(^a\)

| Parameter Name | Value       | Description                                                                 |
|----------------|-------------|------------------------------------------------------------------------------|
| comp_2_damp    | [0.6]       | Initial guess for amplitude of second component as fraction of first component amplitude |
| comp_2_dvel    | [−150, −50, +50, +150] | Initial guess range for velocity of second component; given as difference between velocities of first and second components |
| comp_2_dvdisp  | [+20]       | Initial guess for velocity dispersion of second component                    |

Note.

\(^a\) LZIFU explores all possible combinations of initial guesses of first and second components.

Calzetti (2001). It also assumes a simplistic dust geometry—a screen of dust between the observer and the Balmer-line-emitting regions. Such an assumption might not be true for the complicated kinematics within the post-collision Taffy system.

4. The Kinematics of the Taffy System

As Figure 3 has shown, the spectra are quite complex in the system, and so we present the kinematic results in two ways. First, we present the channel maps of one of the lines (in this case \(\text{H}\alpha\)) to provide a large-scale view of the gas distribution as a function of radial velocity channel. Second, we explore the spatial distribution of the gas associated with different kinematic components, especially those associated with multiple lines.

4.1. \(\text{H}\alpha\) Channel Maps

Figure 6 shows the \(\text{H}\alpha\) emission channel maps integrated over channels of width 70 km s\(^{-1}\). It is well known that the two galaxies are counterrotating (e.g., Vollmer et al. 2012), and the brightest ionized gas in the two systems reflects this. Ionized gas in Taffy-S is seen in the lowest velocity channel (3806–3876 km s\(^{-1}\)) in the northwest on the inside edge of the faint stellar ring and progresses in a southeasterly direction with increasing velocity, eventually showing a major component of emission in the southeast disk that fades away around 4780 km s\(^{-1}\). Taffy-S also has some peculiar kinematics. For example, in the NW part of the disk, faint gas emission is seen over a wide range of velocities along the northern major axis even at the highest velocities. This would not be expected for gas in normal rotation. Taffy-N is even more peculiar. The main centroid of emission from the southeast disk appears at around 4016–4086 km s\(^{-1}\) and progresses steadily toward the northwest, showing the counterrotation. In addition, there is a peculiar region of emission that appears at even lower velocities on the northwest extreme tip of Taffy-N and cannot be part of the normal rotation of the galaxy. Indeed, that structure appears to be part of the bridge, since as velocities increase it becomes more extended and eventually connects to the northwestern region of Taffy-S. In addition to this bridge feature, a second bridge component starts to appear between the galaxies at velocities of 4000 km s\(^{-1}\), and at higher radial velocities it becomes quite strong in the region of the extragalactic H II region. The emission bridges the two galaxies where it joins with emission that potentially is associated with the faint stellar ring in the northeastern part of Taffy-S. The connection between the galaxies disappears at velocities in excess of 4650 km s\(^{-1}\).

The faint stellar ring in the northern half of Taffy-S exhibits some peculiar emission. Features that can be associated with this ring can be seen most clearly appearing from velocities around 4000 km s\(^{-1}\). Moving to higher velocities shows several clumps that appear to follow the ring from NW to SE. These clumps also appear to be surrounded by emission that blends with the emission from the bridge. These features hint that this ring was strongly influenced by the collision and shows a discernible transition from material that is clearly associated with Taffy-S to material clearly associated with the bridge. A similar argument could also apply to Taffy-N, although such a transition is much harder to observe in Taffy-N, which is highly inclined. It is clear that the velocity structure of the gas in both galaxies and the bridge is very complex, and so we will now explore the gas in terms of its spectral profiles—which allows us to more easily separate normal regular rotation in the galaxies from peculiar motions.

4.2. Mapping Two-component Line Profiles in the System

As we have shown previously, there are regions in the full data cube where the emission-line spectra show more than one
component. Similar, double-line profiles were noticed in both the H I (Condon et al. 1993) spectra and spectra taken in the far-IR [C II] and [O I] lines with Herschel (Peterson et al. 2018). Some regions of the bridge also show two components in the CO 1−0 molecular gas observations of Gao et al. (2003). Our observations (consistent with those of Gao et al. 2003; Peterson et al. 2018) show that the double-line profiles in the ionized gas are not just confined to the bridge but are also seen projected against parts of the galaxy disks, especially Taffy-N.

To explore the kinematics further, we performed line fitting in two distinct passes. First, we ran LZIFU, forcing it to fit only one component across the whole system. This worked well in regions where the lines were single valued but produced poor results in regions where the lines were double profiled. The output from LZIFU at this stage was a model data cube built from the model fits, as well as the best-fitting continuum cube. Also included were additional data products, including integrated line maps for each of the lines fitted.

Second, we ran LZIFU again, but this time we required it to fit two components for everything. In this case, the fitting worked well for the case of two components, but we found that it performed poorly in places where the profiles were singular. In this mode, since the software was fitting two separate line profiles, the output included two sets of data products (model line cubes, integrated line and velocity field maps), one set for each component—a low- and high-velocity component.

We were able to interrogate the model results to determine, spaxel by spaxel, the mean, standard deviation, and amplitudes for each of the single- and two-component fits. This allowed us to divide the results into three classes of kinematic behavior:

1. Spectra consistent with a single Gaussian component, $V_1$.
2. Spectra consistent with two components, $V_1$ and $V_2$, at different velocities ($V_1$ represents the lowest-velocity component, and $V_2$ the highest).
3. Spectra consistent with two components with nearly the same mean velocity, but exhibiting both a narrow $V_{1n}$ and broad $V_{2b}$ component.

To qualify as two components, the two Gaussian line centers, $V_1$ and $V_2$, were required to differ by at least 35 km s$^{-1}$, or one-half of the velocity resolution at Hα wavelengths. In practice, the lines were generally farther apart than this and clearly separated. Similarly, for a second component to be considered broad, $V_{2b}$ must have an FWHM of at least 1.5 times that of $V_{1n}$. In a few cases, the LZIFU modeling failed to fit two components to cases where a single line profile would have been more appropriate. In this small number of cases, we inspected the profiles by eye to perform what we considered to be the most reasonable classification.

In general when there was doubt about the classification, we inspected the profiles by eye to confirm the classification. Overall the separation between single and two components seemed reasonable, although we realize that there may be some areas where the distinction is somewhat subjective. Less subjective methods of deciding whether to fit single or multiple
Figure 6. Velocity channel map showing the Hα emission over −540 to +510 km s\(^{-1}\) with respect to the (optically defined) systemic recessional velocity of \(\sim 350 \text{ km s}^{-1}\). The contours have been overlaid on an SDSS i-band image of the Taffy galaxies. The color of the contours, going from white to blue, indicates the intensity of the Hα emission going from low to high intensity. The contour levels are 1, 4, 10, 20, 40, 80, 105, and 135 in units of erg s\(^{-1}\) cm\(^{-2}\)/(70 km s\(^{-1}\)). Note that every contour level is not visible in each panel because the Hα intensities change for each panel. The velocity range shown for each panel is 70 km s\(^{-1}\). The LZIFU emission-line cube was stitched as described in Section 4.2. The scale bar in the top left panel is 1 arcmin × 18 kpc.

Given that the kinematics are quite complex, we start by identifying those regions that may show regular galactic rotation. The simplest kinematics to understand are those of Taffy-S in the low-velocity component of Figure 7(b). Here the increasing isovelocity contours, going from yellow to dark blue, progress regularly in the double-line region, merging smoothly with the \(V_{1n}\) (green underlying color) and single-component \(V_s\) (red underlying color) contours. The velocity dispersion in the low-velocity component in Taffy-S is also low across most of its disk. Concentrating only on the low-velocity component for Taffy-S, it is clear that the velocities and dispersions shown in Figures 7(b) and (c) look like a somewhat warped but regular rotating disk. In contrast, Taffy-S is much more peculiar in the high-velocity component of Figures 7(e) and (f), where the galaxy shows only a small amount of obvious rotation, as well as exhibiting a high velocity executor 4.2. The color of the contours, going from white to blue, indicates the intensity of the Hα emission going from low to high intensity. The contour levels are 1, 4, 10, 20, 40, 80, 105, and 135 in units of erg s\(^{-1}\) cm\(^{-2}\)/(70 km s\(^{-1}\)). Note that every contour level is not visible in each panel because the Hα intensities change for each panel. The velocity range shown for each panel is 70 km s\(^{-1}\). The LZIFU emission-line cube was stitched as described in Section 4.2. The scale bar in the top left panel is 1 arcmin × 18 kpc. The LZIFU emission-line cube was stitched as described in Section 4.2. The scale bar in the top left panel is 1 arcmin × 18 kpc.
dispersion in a large part of the disk. It also has a band of spectra classified as broad-line (Class-3 type; green underlying color) in the nuclear regions.

We now turn our attention to Taffy-N, which has more complex kinematics. This galaxy is quite edge-on and may be expected to show regular rotation along its major axis. Evidence of rotation is seen in the high-velocity component of Taffy-N in the southern part of the disk centered on the dense dust lane and nucleus. Figure 7(e) shows a clear rotation signature where the velocities (starting with the pale-blue contours in the southeast) increase along the major axis (dark-blue contours in the northwest of the inner disk). Although this apparent regular rotation in the high-velocity component is confined to the inner parts of the Taffy-N disk, the increasing trend in velocity shows a reversal toward the northwestern extended disk. This might be interpreted as a turnover in the rotation curve there.

Next, we consider those parts of the velocity field that cannot be considered normal and are most likely caused by the strongly collisional nature of the Taffy pair. We have already pointed out in the discussion of the channel maps that the northwestern part of Taffy-N has a peculiar low-velocity structure that is not part of the normal rotation. This can be seen in Figure 7(b) (low-velocity component), where much of the NW disk of Taffy-N shows very little rotation and also shows high velocity dispersion (Figure 7(c)). Here we see that the gas extends as a finger toward the southwest, where it forms a western bridge with Taffy-S. A second eastern bridge structure is seen associated with the extragalactic H II region, which is strongest in the high-velocity component. The two structures are graphically emphasized in Figure 8. Much of the gas between the two galaxies is seen in the eastern high-velocity bridge component and shows a velocity gradient that extends between the two galaxies along the direction of the radio

![Figure 7](https://example.com/figure7.png)

**Figure 7.** Contour maps, overlaid on an SDSS i-band image of the Taffy system, of the moments of the velocity field. Zero velocity corresponds to an (optically defined) recession velocity of 4350 km s⁻¹. The top and bottom rows correspond to the low- and high-velocity component as shown. From left to right, the column panels show the integrated flux in Hα, radial velocity (with respect to systemic velocity), and velocity dispersion, respectively. The contour levels for the integrated line flux maps are 2000, 3000, 6000, 12,000, 15,000, 20,000, 25,000, and 30,000 in units of erg s⁻¹ cm⁻² km s⁻¹. The contour levels for the velocity maps are −350, −250, −200, −150, −100, 0, 100, 150, 200, 250, and 350 in units of km s⁻¹. The contour levels for the velocity dispersion maps are 50, 70, 90, 130, 160, 190, and 230 in units of km s⁻¹. The red polygon demarcates the boundary where we see two line components in the profiles of the emission lines. The red spaxels outside the double-line boundary indicate spaxels where we see only a single velocity component. The green spaxels indicate spaxels where we see two components but with significantly different widths (i.e., Class-3; we define the three line profile classes in Section 4.2).
5. Excitation of the Ionized Gas in the Taffy System

We next consider the possible excitation mechanisms for the ionized gas within the Taffy system by constructing emission-line diagnostic diagrams based on LZIFU fitting of each spaxel in the data cube (sometimes called BPT or VO diagrams; Baldwin et al. 1981; Veilleux & Osterbrock 1987; Kewley et al. 2001). We construct emission-line diagnostic diagrams using the [O III] $\lambda 5007$/$H\beta$, [N II] $\lambda 6583$/$H\alpha$, [O I] $\lambda 6300$/$H\alpha$, and [S II] $\lambda \lambda 6716, 6731/H\alpha$ ratios. We have significant detections of the [N II] and [S II] lines in most spaxels that fall on the galaxies and bridge, whereas the [O I] $\lambda 6300$ line is detected in fewer spaxels because it is not as strong. We require every emission line used in the diagrams to be detected at the $3\sigma$ level.

In order to look for differences in excitation properties between the low- and high-velocity components where the lines are double, we show line diagnostic diagrams for each of the components separately (panels (a) and (b)) in each of Figures 10–12 for the [N II], [O I], and [S II] plots, respectively. We also show the diagnostic diagrams for the total emission (sum of the two components plus those fitted by a single line) as a third panel (c) in each of the same figures. These diagrams use classifications from Kewley et al. (2006). In all three line diagnostic diagrams, we plot spaxels associated with different spatial regions of the Taffy system. The symbols shown in the excitation diagrams are labeled in each figure based on the regions defined in Figure 3. For example, green circles and crosses represent regions in Taffy-N, while blue circles represent Taffy-S (the nuclear region is distinguished as blue diamonds). The Taffy bridge is shown as red crosses (east bridge) and filled orange circles (west bridge). Because the west bridge contains fewer points than its eastern counterpart, we also present the integrated line ratio for the western bridge as a single larger open orange circle.

The line diagnostic diagrams show that the low- and high-velocity components often behave differently, indicating a difference in their respective excitation mechanisms. For the two galaxies (blue and green symbols), all three sets of line diagnostic diagrams show spaxels mainly distributed within the H II, H II+AGN-composite, or LINER part of the diagnostic diagram, with little hint of any pure AGN component. Much recent work has shown that excitation by shocks can resemble excitation by an AGN. Rich et al. (2011, 2014, 2015) showed that composite line ratios, i.e., H II + AGN, can be due to H II + shocks. This is particularly true for merging galaxies, for example, Rich et al. (2014) showed that merging U/LIRGS can present “composite” optical spectra in the absence of any AGN contribution, with increasing contribution from shocks as the
merger stage progresses from early to late mergers. We will argue below that much of the LINER emission, with the possible exception of the nucleus of Taffy-S, is likely the result of fast shocks exciting the ionized gas in the bridge and in significant parts of the galaxy disks.

We will first concentrate on describing excitation in Taffy-N. We decided to split Taffy-N into two parts, one part covering the main disk of the galaxy and the other part covering the western extension, which appears to connect with the western bridge. The emission from the main disk of Taffy-N (green circles) is largely consistent with emission from HII regions. This is true in all of the diagnostic diagrams in the low, high, and summed components. The western extension of Taffy-N shows a mix of HII and LINER emission. This is especially evident in Figure 11, where the western extension of Taffy-N falls clearly in the LINER area in the low-velocity component, which we have previously noted from the channel maps may be associated with the western bridge.

Next, we consider Taffy-S (blue circles and diamonds). There appears to be a strong mix of HII region and LINER/composite excitation for Taffy-S in all diagnostic diagrams. Despite the fact that the nucleus of Taffy-S does not show evidence for a powerful AGN in the gas excitation diagrams, this does not necessarily preclude an LLAGN being present—especially given the broader line widths discussed previously. Taffy-S’s nucleus (blue diamonds on Figures 10–12) does show some evidence of being a LINER, especially in the low-velocity regime. We find that most of the spaxels located on the nucleus of Taffy-S in the low-velocity ($V_1$) line diagnostic diagrams are in the LINER area in the [O I] and [S II] diagnostic diagrams and close to the AGN line in the [N II] diagram for both the low- and high-velocity components. However, as previously noted, the velocity

Figure 10. [N II] $\lambda$6583 line diagnostic diagrams: (a) for the low-velocity component, (b) for the high-velocity component, and (c) total, i.e., sum of both velocity components and also including line ratios from spaxels that show a single component. Each point here represents the line ratios from a single spaxel. The red crosses, green circles, and blue circles correspond to the eastern bridge, Taffy-N, and Taffy-S, respectively. The blue diamonds are spaxels that fall within the nuclear region of Taffy-S. The green plus signs are spaxels from the western part of Taffy-N. The orange circles are spaxels from the western part of the bridge. The unfilled orange circle is the average of line ratios from all the spaxels within the western bridge. These colors are consistent with those used to denote the corresponding regions in Figures 3 and 8, with the exception of the western part of Taffy-N, which is shown as a magenta polygon in Figure 3. The panels also show line ratios from the MAPPINGS III shock models (Allen et al. 2008) overlaid on the measured line ratios, i.e., colored solid lines. The parameters assumed in the models are $Z = Z_e$ and $B = 5 \mu G$. Along each shock velocity line the points are marked by increasing number density. The classifications are from Kewley et al. (2006). The shaded gray area around the solid black line classifying the HII region excited gas marks the area we used to put a lower limit on the fraction of gas excited by star formation (see Section 6). The width of the shaded area (above the HII classification line) is twice the size of the average error bar in each panel.
difference between the two components is small (<20 km s\(^{-1}\)), with the main difference being in the width of the lines, the \(V_1\) component having a broader width than the \(V_2\) component (see Figure 9). This is consistent with Chandra X-ray observations (Appleton et al. 2015), which showed the possible existence of an LLAGN based on the X-ray hardness ratio.

We divide the bridge into two parts (east and west) as discussed previously. The eastern bridge (red crosses) shows clear evidence of being H II region excited in the low-velocity component. The situation is quite different for the high-velocity \(V_2\) component (“b” panels) in Figures 10(b) and 11(b). Here we observe sets of east-bridge spaxels that deviate strongly from the H II area locus. For example, in the \([\text{O III}]\ \lambda5007/\text{H} \beta\) ratio, the east-bridge points are spread out along the \([\text{N II}]\/\text{H} \alpha\) and \([\text{O I}]\/\text{H} \alpha\) line, extending strongly into the LINER area. The western bridge (orange filled circles) shows a mix of H II region and composite/LINER behavior in all diagrams. This is quite similar to what we find for the western extension of Taffy-N, indicating that they might be excited by the same processes. Because there are so few points from the western bridge, we also plot an average of the entire western bridge as an orange open circle, which is only shown in the “all” components diagram. This falls in the composite/LINER area for the \([\text{N II}]\) and \([\text{O I}]\) diagrams but in the H II area for the \([\text{S II}]\) diagram.

5.1. Evidence for Shocked Gas in the Taffy System

We first explore the possibility that the gas is excited by shocks. We overplot predicted line ratios from shocks on Figures 10–12, taken from the MAPPINGS III library of models (Allen et al. 2008). The model line ratios are plotted for different shock velocities as solid colored lines. Based on the models of Vollmer et al. (2012), we assume that much of the gas in the bridge and throughout the galaxies is close to solar metallicity, since the gas has been stripped from the galaxies or has been excited in situ. For the shock models we therefore assume solar metallicity. The other parameters of the models include the pre-shock gas densities in the range \(0.1 < n \left(\text{cm}^{-3}\right) < 1000\) (stepping by a factor of 10 each time) and an assumed constant magnetic field of \(B = 5 \mu\text{G}\) (this is close to the equipartition magnetic field strength of 8 \(\mu\text{G}\) measured by Condon et al. 1993 through radio continuum measurements). We have plotted only the line ratios for the shock itself while excluding the precursor component of the shock. For the moderate shock velocities that seem compatible with the Taffy excitation velocities, it may be reasonable to ignore the effect of a strong ionizing shock precursor, as we shall discuss later.

For the east-bridge points (red crosses), in the high-velocity component (“b” panels), it is clear that the east-bridge points

![Figure 11](image-url)
fall relatively neatly between the solid lines for shock velocities 175 and 200 km s$^{-1}$ (green and orange lines). For the west-bridge points, in cases where the points deviate from the H II region area they are consistent with the same shock velocity, e.g., the low-velocity components in all diagrams. Thus, the high-velocity east-bridge component and the low-velocity west-bridge component seem consistent with shock excitation for all the line diagnostic diagrams.

The situation is mixed for the disks of galaxies themselves. As Figures 10(b), 11(b), and 12(b) show, there are some points within the disks of both galaxies that fall in the composite region of the diagnostic diagram. Some of these points would be consistent with a mixture of H II region and shocked gas excitation. The nucleus of Taffy-S is an ambiguous case, because it could be excited by shocks in a mild outflow, or may be gas excited by UV emission from a weak LLAGN.

The spreading of the points along lines of constant [O III] $\lambda$5007/H$\beta$ ratio in the high-velocity component in the bridge has a number of possible interpretations if we assume that shocks are involved. First, the spread might imply that the shocks are occurring in an ensemble of gas clouds with different pre-shock densities. Such a picture is consistent with our previous observations of the Taffy bridge (Peterson et al. 2012, 2018), where we have observed gas in many different excited phases, from H I (Condon et al. 1993) to warm molecular gas from the Spitzer IRS, along with the detection of [CI] and [C II] emission (Peterson et al. 2018) and boosted values of [C II]/far-IR and [C II]/PAH ratios. The existence of a highly multiphase (and multidensity) medium is also very consistent with the detection of soft X-ray emission from the bridge. Thus, it might be expected that shocks moving through such a multiphase gas would encounter a range of pre-shock densities—which would spread the points along lines of constant shock velocity—as observed in Figures 10 and 11 especially.

An alternative explanation might be that some of the gas is of lower metallicity. As the models of Allen et al. (2008) show, reducing the metallicity of the shocked gas moves the points in the diagnostic diagrams to the left at roughly constant values of [O III] $\lambda$5007/H$\beta$ ratio. However, if the collisional models of Vollmer et al. (2012) are correct, the gas in the bridge should have come from many different places within the original pre-collisional disks, and deviations of factors of 100 in metallicity in the bridge seem unlikely. We conclude that it is much more likely that we are observing shocks within the bridge and parts of the galaxy disks that encounter clumps of material at different densities.

Figure 12. Same as Figure 10, but using the sum of the [S II] $\lambda$6716, 6731 lines.
What is the effect of ignoring the possible influence of a hot shock precursor in the models? This is an effect where, in high-velocity shocks, the gas in the shock is so strongly heated that UV radiation from the shocked gas ionizes large amounts of pre-shocked gas upstream of the shock. We show in Figure 13 an example of the [N II] line diagnostic diagram, the effect of including the shock and the shock precursor (Allen et al. 2008). As can be seen by comparison with Figure 10, the behavior when we include the shock precursor with velocities $<300 \text{ km s}^{-1}$ is very similar to the case with no shock precursor, which fits our data well. At shock velocities $>300 \text{ km s}^{-1}$ the models including the shock precursor diverge significantly from these data. Similar behavior is noted in the other diagnostic diagrams (not shown). This implies that the shock models between 100 and 300 km s$^{-1}$ fit the bridge data well regardless of whether the precursor is included. We note, however, that in dense gas (e.g., molecular gas known to also be present in the bridge) the velocity at which a shock precursor may become important will be much lower. Therefore, future modeling of the molecular shocks may have to take precursor activity into account.

5.2. Alternatives to Shock Excitation: Diffuse Lyman Continuum Emission Leaking from H II Regions?

In a recent paper by Weilbacher et al. (2018), it was noted that the Antennae galaxies (NGC 4038/39), like the Taffy system, also exhibit significant diffuse ionized gas emission. Although several mechanisms were put forward to explain the emission, including the possibility of shocks, the authors favor an interpretation that much of the diffuse gas is ionized by Lyman continuum photons (hereafter Ly-C) “leaking” from large numbers of H II regions found primarily in the disks of both galaxies, as well as H II regions found within the “overlap region.” Much (but not all) of the diffuse component was found close to massive star formation complexes in the system. Using multicolor HST imaging of the clusters, the authors were able to compare the luminosity of the H$\alpha$ emission associated closely with a cluster with theoretical models of the Ly-C flux from the clusters to determine whether the clusters were “leaking” Ly-C photons into the surrounding gas. It was found that many of the H II regions had nonzero escape fractions of Ly-C UV radiation, especially in the center of NGC 4038 and also in the “overlap” region. Thus, for the Antennae system, it was found that the excess diffuse emission within the system could be explained as gas excited by UV radiation escaping from the clusters.

Could some or all of the extended ionized gas emission we see in the Taffy come from similarly degraded UV light from H II regions in the disks of the Taffy galaxies that might diffuse outward and ionize large parts of the Taffy bridge? We estimate (Table 1) that the amount of H$\alpha$ emission coming from the bridge, after correcting for extinction, is similar to that from Taffy-S and roughly 50% of the emission from Taffy-N. In this case the majority of the escaping photons would have to come from the galaxies themselves, since the SFR in the Taffy bridge is very low. Unfortunately, unlike the case of the Antennae, we do not have multicolor high-resolution images of the individual star clusters, and this means that it is difficult to perform the same kind of test that was applied, for all individual H II regions, by Weilbacher et al. (2018). As a result, we cannot completely rule out a significant contribution to the ionized medium in the Taffy coming from leaky H II regions. Nevertheless, many other lines of evidence already suggest that shocks must be present in the Taffy bridge, and so we prefer the shock explanation for the excitation of the high-velocity component, rather than leaky H II regions. Future observations will be needed to attempt to model the history of star formation in the clusters in the galaxies, which will allow us to estimate the fraction of UV emission that may escape into the surrounding gas. This is beyond the scope of the current paper.

6. Ionized Gas Fractions, SFRs, and Mass in the Ionized Component

Here we describe the method that we employed to estimate a lower limit to the fraction of ionized gas excited by star formation (as opposed to being shock excited) using the [N II] line diagnostic diagram. We use the [N II] line diagnostic diagram since it contains the greatest number of points. We start by defining an effective H II region excitation area. This is
shown as a shaded gray region in the [N II] line diagnostic diagrams of Figure 10. We defined this area simply by "padding" the H II classification line (Kewley et al. 2006) by twice the size of the average error in the y-direction (the y-error being the larger error). We sum up the Hα flux in each spaxel that falls within this region. This flux is divided by the total Hα flux to arrive at the lower limit for the fraction of ionized gas excited by star formation. This process is repeated independently on each velocity component.

The fraction derived this way is a lower limit because the other spaxels, outside of the shaded area (i.e., in the H II+AGN area), will contain emission from gas excited by star formation, and we cannot accurately disentangle the excitation from shocks and star formation (also see text in Section 5.1). The lower limits for the fraction of ionized gas excited by star formation that we derived are 64% and 46% for the lower- and higher-velocity component, respectively. For the purposes of the calculations of SFR (in Section 6.1) and ionized gas mass (in Section 6.2) we estimate the Hα luminosity, coming only from star formation, by

\[ L(\text{H}\alpha) = \frac{0.64 \times L(\text{H}\alpha)_{\text{low}} + 0.46 \times L(\text{H}\alpha)_{\text{high}}}{1} \]

where \( L(\text{H}\alpha)_{\text{low}} \) and \( L(\text{H}\alpha)_{\text{high}} \) are the extinction-corrected total Hα luminosities in the low- and high-velocity components, respectively.

This gives us an extinction-corrected value of \( L(\text{H}\alpha)_{\text{SF}} = (4.99 \pm 0.54) \times 10^{41} \text{ erg s}^{-1} \) for the lower limit to the Hα luminosity resulting from star formation for the Taffy system.

### 6.1. SFR Estimate from Hα Luminosity

Using the following relation from Kennicutt (1998), we estimate the SFR in the entire Taffy system and the bridge region (using the entire region defined as the bridge in Figure 3), including Hα emission from the extragalactic H II region:

\[ \psi[M_\odot \text{ yr}^{-1}] = 7.9 \times 10^{-42} L(\text{H}\alpha)[\text{ erg s}^{-1}] \]

We obtain

\[ L(\text{H}\alpha)_{\text{SF}} = (4.99 \pm 0.54) \times 10^{41} \text{ erg s}^{-1} \]

and

\[ L(\text{H}\alpha)_{\text{SF, bridge}} = (1.02 \pm 0.14) \times 10^{41} \text{ erg s}^{-1} \]

for the extinction-corrected Hα luminosity coming from star formation for the entire Taffy system and bridge, respectively, using the method described previously. This translates to SFRs of \( 3.94 \pm 1.0 M_\odot \text{ yr}^{-1} \) and \( 0.81 \pm 0.22 M_\odot \text{ yr}^{-1} \), respectively, for the entire Taffy system and the bridge. These SFRs agree well with our previous estimates derived from UV–far-IR SED fitting (Appleton et al. 2015) of \( 3.65 \pm 0.03 M_\odot \text{ yr}^{-1} \) and \( 0.69 \pm 0.06 M_\odot \text{ yr}^{-1} \), respectively, for the entire system and the bridge.

Interestingly, recent observations with the Atacama Large Millimeter Array (ALMA) show dense filaments of molecular gas, in the Taffy bridge, with little star formation in them. These ALMA observations and the overall star formation properties will be discussed in a future paper (P. N. Appleton et al. 2019, in preparation).

### 6.2. Ionized Gas Mass

The mass of ionized gas in the Taffy system and in the bridge assuming case B recombination (Macchetto et al. 1996; Kulkarni et al. 2014) is given by

\[ M_{\text{ion}} = 2.33 \times 10^{3} \left( \frac{L(\text{H}\alpha)}{10^{29}} \right) \left( \frac{10^{3}}{n_{e}} \right) M_\odot \]

where \( L(\text{H}\alpha) \) is the extinction-corrected Hα luminosity in units of erg s\(^{-1}\) and \( n_{e} \) is the electron density. For the bridge, from the lower limits to the ionized gas fractions from star formation (see above) we obtained \( L(\text{H}\alpha)_{\text{SF, bridge}} = (1.02 \pm 0.14) \times 10^{41} \text{ erg s}^{-1} \). We also have \( n_{e} = 200 \text{ cm}^{-3} \), based on the ratio of the [S II] lines.

This then gives us \( M_{\text{ion}} = (1.19 \pm 0.22) \times 10^{5} M_\odot \). This calculation is uncertain because the Hα emission originates from two different processes—H II regions and very likely shocks. However, it does show that the ionized gas mass is an insignificant fraction (~0.2%) of the total mass of gas in the bridge (~7 \times 10^{9} M_\odot), made up of a mix of H II and H3. This is in agreement with the very low ionized gas fraction responsible for exciting the [C II] 157.7 μm far-IR cooling line in the bridge (Peterson et al. 2018) determined from the upper limit to the detection of [N II] 206 μm in the bridge. For the Taffy system as a whole, using the extinction-corrected Hα luminosity coming from star formation, \( L(\text{H}\alpha)_{\text{SF}} = (4.99 \pm 0.54) \times 10^{41} \text{ erg s}^{-1} \), we get \( M_{\text{ion}} = (5.8 \pm 1.0) \times 10^{5} M_\odot \) for the mass of ionized gas in the entire Taffy system. Again, this is an insignificant fraction (~0.8%) of the total gas mass in the Taffy system.

### 6.3. Post-starburst Populations

We detect Hβ absorption lines within many spaxels on the galaxies. The spectra of post-starburst galaxies are known to contain strong Balmer absorption lines owing to their stellar populations being dominated by A-type stars. Evidence of a post-starburst population is not uncommon in merging galaxies (e.g., Zabludoff et al. 1996; Yang et al. 2004, 2008), but attempts to measure the age of the stellar population are difficult, especially when only H β is observed (see, e.g., Worthey & Ottaviani 1997). Since our spectral coverage did not include other post-starburst indices, we can only provide preliminary results here. Further observations using full UV–optical SEDs, better absorption-line indices, and detailed modeling (e.g., French et al. 2016) will be needed to obtain a better estimate for the age of the population that is responsible for the Hβ absorption.

We measured the EW of the Hβ absorption line for each spaxel that contained either the galaxies or the bridge. The EW was measured using

\[ W(\text{H}\beta) \, [\text{Å}] = \frac{\int_{\text{line}} f_{\lambda} d\lambda}{\langle f_{\lambda,\text{cont}} \rangle}, \]

where the integral is done over the continuum-subtracted absorption-line fit and \( \langle f_{\lambda,\text{cont}} \rangle \) is the average continuum value measured on either side of the Hβ line. Lzifu provides as output the fit to the stellar continuum and the nebular emission lines separately. We use the continuum fit cube to refit a Gaussian absorption line to the region centered on the Hβ absorption. The parameters from this fit then give the area within the absorption line, and the average continuum is measured in a band of ~10 spectral elements on both sides of the line.

Figure 4(b) shows the measured EW map for the Taffy system. Relatively deep Hβ absorption (W(Hβ) > 10) occurs where there is little Hα emission. Two regions of high W(Hβ)
lie in the northern and southern parts of the faint stellar ring that surrounds Taffy-S. Another region, in the southeastern part of Taffy-N, does extend into regions where there is some star formation and older stellar populations are probably present, with the deepest absorption lying outside of the main star formation disk.

It is generally true that values of $W(H\beta)$ in the range of 10–20 $W(H\beta)$ imply stellar evolutionary ages for the post-starburst populations of several hundred megayears, and this would imply that there is no connection between the post-starburst population and the current collision between the two galaxies (e.g., Worthey & Ottaviani 1997). From dynamical arguments it has been argued that the collision between the two Taffy galaxies is quite recent (approximately 25–30 Myr; Vollmer et al. 2012), and so the fact that the outer ring of Taffy-S shows an old population would lead to an apparent problem, if the stellar population of the ring were created in the collision. However, one solution might be that the stellar population of the ring is a stellar density wave containing a much older pre-collisional system that had undergone star formation a long time in the past. This is also the conclusion reached by Jarrett et al. (1999), who argue that the ring consists of stars from an old disk population from the pre-collision disk.

The fact that both galaxies contain evidence of post-starburst activity might imply that the galaxies underwent a high-speed encounter in the more distant past that triggered star formation but did not lead to immediate merger. If that is the case, then we are probably currently witnessing the second (probably final) collision before full merger. Broader wavelength coverage in the blue, to detect more absorption lines and better characterize the age of the post-starburst population, will be needed to test this idea.

7. Conclusions

Using visible IFU data, from the VIRUS-P instrument on the 2.7 m telescope at McDonald Observatory, for the Taffy system, we have shown the results summarized below.

1. We detect widespread ionized gas within the disks of the Taffy galaxies and the bridge that exhibit very disturbed kinematics, including many regions with double-line profiles and emission regions that do not follow regular rotation. Although both galaxies show velocity components that approximate gas rotation around their centers, both galaxies also show peculiar motions, often associated with gas that extends into the bridge between them. The gas associated with Taffy-N (UGC 12915) contains a major kinematic component that does not take part in regular rotation but exhibits a high velocity dispersion and forms a narrow western bridge to Taffy-S (UGC 12914). Taffy-S, although showing the kinematics of a likely tidally warped but regular (counter)rotating disk, also contains a peculiar kinematic component that is associated with a second ionized gas bridge. This eastern bridge component, which extends from Taffy-N, through the region containing the extragalactic H II region, and eventually connecting with Taffy-S, is more closely associated with the molecular bridge seen previously to extend between the galaxies. On the other hand, the western bridge, which is much fainter, appears to be kinematically linked with the western extension of Taffy-N (and the western part of Taffy-S) and shows a mix of H II region and composite/LINER excitation unlike the main disk of Taffy-N, which is largely consistent with H II region excitation.

2. An analysis of the excitation of the ionized gas through diagnostic line ratios shows that a significant fraction of the emission shows a mix of H II region and LINER-type emission, especially in the areas where we can clearly distinguish two velocity components. The LINER-type emission is especially dominant in the high-velocity component in the east-bridge region, but also over significant portions of both galaxies. We observe emission-line ratios in the high-velocity component, for the east bridge and the low-velocity component of the west bridge, that are consistent with the gas being excited by shocks with velocities of $\sim$175–200 km s$^{-1}$ and a range of pre-shock densities. Such evidence for shocks permeating clouds of varying density is consistent with previous observations of the bridge by Spitler and Herschel, where strong mid- and far-IR cooling lines are detected from warm molecular and diffuse atomic gas heated by turbulence. While we cannot rule out a contribution to the diffuse ionized emission from Lyman continuum photons leaking from young H II regions embedded in the disks of both galaxies, the weight of evidence, from previous multiwavelength observations of the bridge, suggests that shocks are a more likely explanation for the LINER-type emission-line ratios seen in the bridge and in parts of the disks of both galaxies. Given the violence of the collision based on previous numerical models of head-on collisions that suggest that the bridge is still in a highly disturbed state, such shocks may not be unexpected.

3. Strong Balmer absorption lines ($10 < W(H\beta) [\AA] < 15$) are observed in parts of the ring associated with Taffy-S (UGC 12914), as well as the southeast portion of the edge-on disk of Taffy-N (UGC 12915). The absorption lines are strongest in regions where the ionized and molecular gas distributions are weak, suggesting that parts of the Taffy system have experienced a burst of star formation in the past, perhaps from a previous close passage of the two galaxies. If so, it is possible that the current collision may be a second, more dissipative collision, which will likely lead to a merger in the near future. Further observations, with a broader blue wavelength coverage than the current observations, will be necessary to better determine the age of the post-starburst populations in both galaxies.

4. Although there has been only weak evidence in the past for the nucleus of Taffy-S (UGC 12914) containing an LLAGN from X-ray properties (Appleton et al. 2015), we detect line widths in the nucleus as large as 320 km s$^{-1}$. These line widths are typical of narrow-line Seyfert galaxies. The broad lines appear to extend over 6–10$''$ along the minor axis. The excitation properties of the nuclear gas are consistent with LINER emission, but without higher spatial and velocity resolution data we cannot determine whether the Taffy-S nucleus hosts a weak shocked, highly confined outflow from a nuclear starburst or is excited by UV radiation from an LLAGN.

5. We provide evidence, supporting much previous work, that the Taffy system has atypically low SFRs for a system having recently undergone a recent major merger. We find
SFRs of $3.94 \pm 1.0 \, M_\odot \, yr^{-1}$ and $0.81 \pm 0.22 \, M_\odot \, yr^{-1}$ in the entire Taffy system and the bridge (including the extragalactic H II region), respectively. Low SFRs in this recent post-collisional remnant may result from the highly disturbed nature of the gas in the galaxies and bridge.

B.A.J. would like to thank the Visiting Graduate Fellowship Program at IPAC/Caltech for six months of support toward work performed for this paper. B.A.J. is also grateful to Drs. Rogier Windhorst and Rolf Jansen for helpful discussions on work done in this paper. The authors thank the anonymous referee for their helpful review and suggestions. This research has made use of NASA’s Astrophysics Data System. This research has made use of the Python programming language along with the Numpy, Scipy, and Matplotlib packages. This research has also made use of Astropy, a community-developed core Python package for Astronomy (The Astropy Collaboration 2018).

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