The Mirror Universes

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Abstract

In this paper we investigate the structure of the Mirror Universes. The two universes are coupled with transformation $t \rightarrow -t$. It is shown that for Planck scale i.e. for $t \cong T_{\text{Planck}} = \left( \frac{\hbar G}{c^4} \right)^{1/2}$ the oscillations of temperature of the universes are observed. For the Mirror Universes the temperature fields are shifted in phase.

Key words: Gravity, universe temperature, oscillation of temperature.
1 Introduction

The physical phenomenon of gravity, described to a high degree of accuracy by Isaac Newton’s mathematics in 1687, has played a key role in scientific understanding. However, in 1915 Einstein created a major revolution in our scientific worldview. According to Einstein theory gravity plays a unique role in physics for several reasons [1]. Most particularly, these are: (i) gravity is the only physical quality that influences causal relationships between space-time events, (ii) Gravitation force has no local reality, as it can be eliminated by a change in space-time coordinates; instead gravitational tidal effects provide a curvature for the very space-time in which all other particles and forces are contained. It follows from this that gravity cannot be regarded as some kind of emergent phenomenon secondary to other physical effects, but is a fundamental component of physical reality.

According to modern accepted physical pictures, reality is rooted in three-dimensional space and a one-dimensional time, combined together into four dimensional space-time. As was stated in Penrose-Hameroff model conscious events are embedded at the Planck scale, and is gravity dependent [2].

In this paper we study the influence of gravity on the thermal phenomena at two-dimensional blisters (1D space, 1D time) in Planck scale. We will consider the Mirror Universes, i.e. two set of the Universes coupled by the transformation $t \rightarrow -t$. It will be shown that the temperature of the Universes are shifted in phase. For short times i.e. for Planck scale, the oscillations of the Universe temperature are predicted.

In this paper we follow of idea of the repulsive gravity as the source of the space-time expansion. We will study the influence of the repulsive gravity ($G < 0$) on the temperature field in the universe. To that aim we will apply the quantum hyperbolic heat transfer equation (QHT) formulated in our earlier papers [3, 4].

When substitution $G \rightarrow -G$ is performed in QHT the Schrödinger type equation is obtained for the temperature field. In this paper the solution of
QHT will be obtained. The resulting temperature is a complex function of space and time. We argue that because of the anthropic limitation of the observers it is quite reasonable to assume \( \text{Im}T = 0 \). From this anthropic condition the discretization of the space radius \( R = [(4N\pi + 3\pi)L_P]^{1/2}(ct)^{1/2} \), velocity of expansion \( v = (\pi/4)^{1/2}((N + 3/4)/M)^{1/2}c \) and acceleration of expansion \( a = -\frac{1}{2}(\pi/4)^{1/2}((N + 3/4)^{1/2}/M^{3/2})(c^7/(\hbar G))^{1/2} \) are obtained.

2 The model

In papers [3, 4] the quantum heat transport equation in a Planck Era was formulated:

\[
\tau \frac{\partial^2 T}{\partial t^2} + \frac{\partial T}{\partial t} = \frac{\hbar}{M_P} \nabla^2 T. \tag{1}
\]

In equation (1) \( \tau = ((\hbar G)/c^5)^{1/2} \) is the relaxation time, \( M_P = ((\hbar c)/G)^{1/2} \) is the mass of the Planck particle, \( \hbar, c \) are the Planck constant and light velocity respectively and \( G \) is the gravitational constant. The crucial role played by gravity (represented by \( G \) in formula (1)) in a Planck Era was investigated in paper [4].

For a long time the question whether, or not the fundamental constant of nature \( G \) vary with time has been a question of considerable interest. Since P. A. M. Dirac [5] suggested that the gravitational force may be weakening with the expansion of the Universe, a variable \( G \) is expected in theories such as the Brans-Dicke scalar-tensor theory and its extension [6, 7]. Recently the problem of the varying \( G \) received renewed attention in the context of extended inflation cosmology [8].

It is now known, that the spin of a field (electromagnetic, gravity) is related to the nature of the force: fields with odd-integer spins can produce both attractive and repulsive forces; those with even-integer spins such as scalar and tensor fields produce a purely attractive force. Maxwell’s electro-dynamics, for instance can be described as a spin one field. The force from
this field is attractive between oppositely charged particles and repulsive between similarly charged particles.

The integer spin particles in gravity theory are like the graviton, mediators of forces and would generate the new effects. Both the graviscalar and the graviphoton are expected to have the rest mass and so their range will be finite rather than infinite. Moreover, the graviscalar will produce only attraction, whereas the graviphoton effect will depend on whether the interacting particles are alike or different. Between matter and matter (or antimatter and antimatter) the graviphoton will produce repulsion. The existence of repulsive gravity forces can to some extent explains the early expansion of the Universe [5].

In this paper we will describe the influence of the repulsion gravity on the quantum thermal processes in the universe. To that aim we put in equation (1) \( G \rightarrow -G \). In that case the new equation is obtained, viz.

\[
\frac{i}{\hbar} \partial T \partial t = \left( \frac{\hbar^3 |G|}{c^5} \right)^{1/2} \partial^2 T \partial t^2 - \left( \frac{\hbar^3 |G|}{c^5} \right)^{1/2} \nabla^2 T.
\] (2)

For the investigation of the structure of equation (2) we put:

\[
\frac{\hbar^2}{2m} = \left( \frac{\hbar^3 |G|}{c^5} \right)^{1/2}
\] (3)

and obtains

\[
m = \frac{1}{2} M_P
\]

with new form of the equation (2)

\[
\frac{i}{\hbar} \partial T \partial t = \left( \frac{\hbar^3 |G|}{c^5} \right)^{1/2} \partial^2 T \partial t^2 - \frac{\hbar^2}{2m} \nabla^2 T.
\] (4)

Equation (4) is the quantum telegraph equation discussed in paper [4]. To clarify the physical nature of the solution of equation (4) we will discuss
the diffusion approximation, i.e. we omit the second time derivative in equation (4) and obtain

\[ i\hbar \frac{\partial T}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 T. \quad (5) \]

Equation (5) is the Schrödinger type equation for the temperature field in a universe with \( G < 0 \).

Both equation (5) and diffusion equation:

\[ \frac{\partial T}{\partial t} = \frac{\hbar^2}{2m} \nabla^2 T \quad (6) \]

are parabolic and require the same boundary and initial conditions in order to be “well posed”.

The diffusion equation (6) has the propagator [10]:

\[ T_D(\vec{R}, \Theta) = \frac{1}{(4\pi D\Theta)^{3/2}} \exp \left[ -\frac{R^2}{2\pi \hbar \Theta} \right], \quad (7) \]

where

\[ \vec{R} = \vec{r} - \vec{r}', \quad \Theta = t - t'. \]

For equation (5) the propagator is:

\[ T_s(\vec{R}, \Theta) = \left( \frac{M_P}{2\pi \hbar \Theta} \right)^{3/2} \exp \left[ -\frac{3\pi i}{4} \right] \cdot \exp \left[ \frac{iM_P R^2}{2\pi \hbar \Theta} \right] \quad (8) \]

with initial condition \( T_s(\vec{R}, 0) = \delta(\vec{R}) \).

In equation (5) \( T_s(\vec{R}, \Theta) \) is the complex function of \( \vec{R} \) and \( \Theta \). For anthropic observers only the real part of \( T \) is detectable, so in our description of universe we put:

\[ \text{Im} T(\vec{R}, \Theta) = 0. \quad (9) \]

The condition (9) can be written as (bearing in mind formula (8)):

\[ \sin \left[ -\frac{3\pi}{4} + \left( \frac{R}{L_P} \right)^2 \frac{1}{4\Theta} \right] = 0, \quad (10) \]
where $L_P = \tau_P c$ and $\tilde{\Theta} = \Theta / \tau_P$. Formula (10) describes the discretization of $R$

$$R_N = [(4N \pi + 3\pi)L_P]^{1/2}(tc)^{1/2},$$

$$N = 0, 1, 2, 3 \ldots$$

In fact from formula (11) the Hubble law can be derived

$$\frac{\dot{R}_N}{R_N} = H = \frac{1}{2t^2}, \quad \text{independent of } N.$$ (12)

In the subsequent we will consider $R$ (11), as the space-time radius of the $N-$ universe with “atomic unit” of space $L_P$.

It is well known that idea of discrete structure of time can be applied to the “flow” of time. The idea that time has “atomic” structure or is not infinitely divisible, has only recently come to the fore as a daring and sophisticated hypothetical concomitant of recent investigations in the physics elementary particles and astrophysics. Yet in the Middle Ages the atomicity of time was maintained by various thinkers, notably by Maimonides [11]. In the most celebrated of his works: The Guide for perplexed he wrote: Time is composed of time-atoms, i.e. of many parts, which on account of their short duration cannot be divided. The theory of Maimonides was also held by Descartes [12].

The shortest unit of time, atom of time is named chronon [13]. Modern speculations concerning the chronon have often be related to the idea of the smallest natural length is $L_P$. If this is divided by velocity of light it gives the Planck time $\tau_P = 10^{-43}$ s, i.e. the chronon is equal $\tau_P$. In that case the time $t$ can be defined as

$$t = M\tau_P, \quad M = 0, 1, 2, \ldots$$ (13)

Considering formulae (8) and (13) the space-time radius can be written as

$$R(M, N) = (\pi)^{1/2}M^{1/2}\left(N + \frac{3}{4}\right)^{1/2}L_P, \quad M, N = 0, 1, 2, 3, \ldots$$ (14)
Formula (14) describes the discrete structure of space-time. As the $R(M, N)$ is time dependent, we can calculate the velocity, $v = \frac{dR}{dt}$, i.e. the velocity of the expansion of space-time

$$v = \left(\frac{\pi}{4}\right)^{1/2} \left(\frac{N + 3/4}{M}\right)^{1/2} c,$$  \hspace{1cm} (15)

where $c$ is the light velocity. We define the acceleration of the expansion of the space-time

$$a = \frac{dv}{dt} = -\frac{1}{2} \left(\frac{\pi}{4}\right)^{1/2} \left(\frac{N + 3/4}{M^{3/2}}\right)^{1/2} \frac{c}{\tau_P}.$$  \hspace{1cm} (16)

Considering formula (15) it is quite natural to define Planck acceleration:

$$A_P = \frac{c}{\tau_P} = \left(\frac{c^7}{\hbar G}\right)^{1/2} = 10^{51} \text{ ms}^{-2}$$  \hspace{1cm} (17)

and formula (16) can be written as

$$a = -\frac{1}{2} \left(\frac{\pi}{4}\right)^{1/2} \left(\frac{N + 3/4}{M^{3/2}}\right)^{1/2} \left(\frac{c^7}{\hbar G}\right)^{1/2}.$$  \hspace{1cm} (18)

In table I the numerical values for $R, v$ and $a$ are presented. It is quite interesting that for $N, M \rightarrow \infty$ the expansion velocity

$v < c$ in complete accord with relativistic description. Moreover for $N, M \gg 1$ the $v$ is relatively constant $v \sim 0.88 c$. From formulae (11) and (15) the Hubble parameter $H$, and the age of our Universe can be calculated

$$v = HR, \quad H = \frac{1}{2M\tau_P} = 5 \cdot 10^{-18} \text{ s}^{-1},$$

$$T = 2M\tau_P = 2 \cdot 10^{17} \text{ s} \sim 10^{10} \text{ years},$$  \hspace{1cm} (19)

which is in quite good agreement with recent measurement [15, 16, 17].
Table I: Radius, velocity and acceleration for N, M-universes

| N, M | R [m]          | v [m/s] | a [m/s²]       |
|------|----------------|---------|----------------|
| $10^{20}$ | $1.77 \cdot 10^{-15}$ | $2.66 \cdot 10^8$ | $-1.32 \cdot 10^{31}$ |
| $10^{60}$ | $1.77 \cdot 10^{25}$ | $2.66 \cdot 10^8$ | $-1.32 \cdot 10^{-10}$ |
| $10^{80}$ | $1.77 \cdot 10^{45}$ | $2.66 \cdot 10^8$ | $-1.32 \cdot 10^{-29}$ |

(*) Spergel D. N. at al. [16];
(++) Anderson J. D. at al. [17]; Radio metric data from Pionier 10/11, Galileo and Ulysses Data indicate and apparent anomalous, constant, acceleration acting on the spacecraft with a magnitude $\sim 8.5 \cdot 10^{-10}$ m/s².

In Fig. 1a,b we present the behaviour of $T_S\left(\frac{R}{R_{\text{Planck}}}, \frac{t}{T_{\text{Planck}}}\right)$ for positive as well negative values of $R$ and $T$. For small values of $R$, $t$ (in comparison to $R_{\text{Planck}}$ and $T_{\text{Planck}}$) one can observe the vibrating structure of $T_S$ (formula (8)) for positive as well negative, $R$ and $t$ values. For the mirror universe $t \to -t$, $R \to -R$ (in 1D case) the "phase shift" is observed.
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Figure 1: The temperature fields for Planck scales, (a) for $\frac{t}{T_{\text{Planck}}} (-100, 100)$
Figure 2: The temperature fields for Planck scales, (b) for $\frac{t}{T_{\text{Planck}}}(-1000, 1000)$