ON THE RATE OF CONVERGENCE OF A NUMERICAL SCHEME FOR FRACTIONAL CONSERVATION LAWS WITH NOISE

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Abstract. We consider a semi-discrete finite volume scheme for a degenerate fractional conservation laws driven by a cylindrical Wiener process. To describe the problem, we assume that \((\Omega, \mathbb{F}, \{\mathbb{F}_t\}_{t\geq 0})\) is a filtered probability space satisfying the usual hypothesis i.e., \(\{\mathbb{F}_t\}_{t\geq 0}\) is a right-continuous filtration such that \(\mathbb{F}_0\) contains all the \(\mathbb{P}\)-null subsets of \((\Omega, \mathbb{F})\). We are interested in numerical approximations of \(L^2(\mathbb{R}^d)\)-valued predictable processes \(u(t, \cdot)\) which satisfy the following Cauchy problem

\[
\begin{cases}
du(t, x) - \text{div} (u(t, x)) \, dt + \mathcal{L}_\lambda [A(u(t, \cdot))](x) \, dt = \sigma(u(t, x)) \, dW(t), & (x, t) \in \Pi_T, \\
u(0, x) = u_0(x), & x \in \mathbb{R}^d,
\end{cases}
\]

where \(\Pi_T = \mathbb{R}^d \times (0, T)\) with \(T > 0\) fixed. Here \(u_0 : \mathbb{R}^d \rightarrow \mathbb{R}\) is a given initial function, \(f : \mathbb{R} \rightarrow \mathbb{R}^d\), \(A : \mathbb{R} \rightarrow \mathbb{R}\) are given (sufficiently smooth) nonlinear functions. Moreover, \(\mathcal{L}_\lambda[u]\) denotes the classical fractional Laplace operator of order \(\lambda \in (0, 1)\), and defined as

\[
\mathcal{L}_\lambda[\psi](x) := d\lambda \text{P.V.} \int_{|z| > 0} \frac{\psi(x) - \psi(x + z)}{|z|^{d+2\lambda}} \, dz,
\]

for some constants \(d\lambda > 0\), and a sufficiently smooth function \(\psi\). The nonlinearity \(A\) is allowed to be zero on an interval, thus (1.1) is a strongly degenerate fractional problem. Furthermore, \(W(t)\) is a cylindrical Wiener process given by \(W(t) = \sum_{k \geq 1} g_k \beta_k(t)\), where \(\beta_k\) are mutually independent real valued standard Wiener processes, and \((g_k)_{k \geq 1}\) is a complete orthonormal system in a separable Hilbert space \(\mathbb{H}\). Finally, we consider the mapping \(\sigma(u) : \mathbb{H} \rightarrow L^2(\mathbb{R}^d)\), for each \(w\) in \(L^2(\mathbb{R}^d)\), defined by \(\sigma(u)g_k = h_k \langle u, \cdot \rangle\).

For our purpose, we assume that \(h_k\) is Lipschitz-continuous and define \(G^2(r) := \sum_{k \geq 1} h_k^2(r)\).

1.1. Review of the existing literature. First observe that, for the case \(A = \sigma = 0\), the equation (1.1) represents a well-known conservation law in \(\mathbb{R}^d\). We refer to the pioneer papers by Kružkov [33] and Vol’pert [36] for existence and uniqueness results related to scalar conservation laws. Numerical schemes for deterministic hyperbolic conservation laws have been studied in [34], [23], [16] as well as others. Well-posedness theory for the deterministic counterpart of (1.1) (i.e., the case \(\sigma = 0\)) has been well studied in literature, starting with the work by Alibaud [1], and Cifani & Jakobsen [15]. For the convergence/rate of convergence of numerical schemes for deterministic degenerate conservation laws one can refer to works by Cifani & Jakobsen [14], Droniou [21], Karlsen et al. [27], and the references therein.

On the other hand, well-posedness theory for stochastic conservation laws (i.e., the case \(A = 0\)) has been established by several authors, see [2, 4, 13, 17]. Regarding the convergence/rate of convergence of numerical schemes for stochastic conservation laws, we mention the paper by Kroker and Rohde [32], authors of this paper and Majee [29], and a series of papers by Bauzet et al. [5], [6], [7]. Moreover,

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Degenerate stochastic equations has been studied by many authors, see [3, 18, 28]. For the stochastic non-local equation (1.1), existence & uniqueness results have been recently developed, by exploiting a new technical framework for the proof of uniqueness, in [8] (linear diffusion case), and [9] (degenerate diffusion case). Moreover, an explicit continuous dependence estimate on the nonlinearities and a rate of convergence estimate for the stochastic vanishing viscosity method was also established in [8, 9]. Given the fact that, a plethora of phenomena in physics and finance are modeled by equations of the form (1.1) (e.g., fluid flows in porous media, and pricing derivative securities in financial markets), there is a pressing need for efficient numerical schemes (such as finite volume schemes) to approximate the underlying problem (1.1). Note that, the challenges in dealing with numerical schemes for SPDEs like (1.1) are manifold, mainly due to the presence of the fractional operator & multiplicative noise term in (1.1). Indeed, one needs to borrow and merge ideas from numerical methods for SDEs and approximation methods for the underlying deterministic problems. This is of course easier said than done since one needs to successfully capture the noise-noise interaction as well. In the realm of well-posedness theory for SPDEs, noise-noise interaction term plays a pivotal role, for details see [8, 24, 11, 12, 10, 30, 31].

1.2. Aims and scope of this paper. Due to the non availability of the explicit solution for stochastic balance laws, we intend to construct a numerical scheme for approximating the solution of (1.1). Note that, although there are large number of papers exploring convergence results for stochastic conservation laws, a rigorous theoretical study for stochastic fractional conservation laws is in a state of infancy. In fact, the specific question about deriving the convergence rate for the approximate solutions to stochastic non-local problems is virtually untouched. For stochastic conservation laws, the convergence rate for monotone methods is known to be $\Delta x^{1/2}$, $\Delta x$ being the discretization parameter. Keeping in mind that the rate of convergence estimate intimately related to the so-called Kato’s type of inequality, the main difficulties in dealing with the specific stochastic non-local problem (1.1) can be summarized as follows: 

1. For the stochastic non-local equation, by virtue of Itô’s formula, one requires to work with a smooth approximations of the usual Kružkov’s entropies.

2. The numerical analysis of the deterministic counterpart of (1.1) heavily relies on the following (Kato’s type of) inequality 

$$
\text{sign}(u(x) - v(y)) \left( \mathcal{L}_\lambda [A(v)](y) - \mathcal{L}_\lambda [A(u)](x) \right) \leq \mathcal{L}_\lambda \left[ |A(u) - A(v)| \right](x,y),
$$

and the above inequality (1.2) does not hold for approximations of the usual Kružkov’s entropies.

In view of the above incompatibility, to establish stochastic Kato’s inequality (even for a semi-discrete scheme!) one has to look beyond the traditional approach of proving convergence of numerical schemes. Indeed, the main technical achievement of this paper stems from successful demonstration of the stochastic Kato’s inequality. The method of proof is new and requires significant changes in computing hierarchical limits with respect to various parameters involved in the proof leading up to Kato’s inequality. In a nutshell, we send the parameter $\delta \downarrow 0$ before sending the parameter $l \downarrow 0$, see Section 4 for a complete description of the main ingredients of our method. Inevitably, above changes in hierarchical limits has effects on all the terms involved in the entropy inequality, and appropriate changes are required to deal with them.

To summarize, we consider a semi-discrete finite difference scheme and show that the expected value of the $L^1$-difference of the approximate numerical solution and the unique entropy solution to (1.1) converges at some rate depending on $\lambda$, which is in accordance with the rate for the deterministic fractional conservation law [14]. Finally, let us mention that we do not know how to derive rate of convergence estimates for a fully-discrete scheme for (1.1), but we strongly believe that the technical achievement of this paper is absolutely essential to establish such a result.

We organize the details as follows. We first briefly recall the entropy framework, introduce the semi-discrete finite volume scheme, and state the main result towards the end of Section 2. Further in Section 3 we prove a priori estimates for the approximate numerical solutions. Section 4 deals with the proof of the main theorem. Finally, we conclude by presenting some numerical results in Section 5.

2. Mathematical Framework and Statement of the Main Result

Throughout this paper, by the letter $C$, we denote various generic constants which may change from line to line. Given a separable Banach space $Z$, let us denote by $N^Z_{2} (0, T; Z)$ the space of square integrable
predictable $Z$-valued processes (cf. [35] p.28 for example). The Euclidean norm on $\mathbb{R}^d$ is denoted by $|| \cdot ||$ and the BV semi-norm is by $| \cdot |_{BV(\mathbb{R}^d)}$. We remark that the space $BV(\mathbb{R}^d)$ consists of functions with bounded variation on $\mathbb{R}^d$, endowed with the semi-norm $|u|_{BV(\mathbb{R}^d)} = TV_s(u)$, where $TV_s(u)$ is the total variation of $u$ defined on $\mathbb{R}^d$. We also denote by $\text{sign}_0(x) = \frac{x}{|x|}$ if $x \neq 0$, and 0 otherwise.

In the rest of the paper, we consider the following assumptions:

A.1 The initial function $u_0$ is a deterministic function in $L^2(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d) \cap BV(\mathbb{R}^d)$.

A.2 The flux function $f = (f_1, f_2, \ldots, f_d) : \mathbb{R} \rightarrow \mathbb{R}^d$ is a Lipschitz continuous function with $f_k(0) = 0$, for all $1 \leq k \leq d$.

A.3 The function $A : \mathbb{R} \rightarrow \mathbb{R}$ is a non decreasing Lipschitz continuous function with $A(0) = 0$.

A.4 There exists $M > 0$ such that $\sigma(u) = 0$, for all $|u| > M$ and $h_k(0) = 0$, for all $k \geq 1$. Moreover, there exists a constant $K > 0$ such that

$$
\sum_{k \geq 1} |h_k(u) - h_k(v)|^2 \leq K|u - v|^2,
$$

which implies $G^2(u) = \sum_{k \geq 1} h_k^2(u) \leq K|u|^2$, for all $u, v \in \mathbb{R}$.

2.1. Stochastic Entropy Formulation. It is well-known that weak solutions to (1.1) need not be unique. Consequently, an entropy admissibility condition must be imposed to single out the physically correct solution. To describe the entropy framework for (1.1), we need to first split the non-local operator $\mathcal{L}_A$ into two terms: for each $r > 0$, we write $\mathcal{L}_A[\psi] := \mathcal{L}_{A,r}[\psi] + \mathcal{L}_A^s[\psi]$, where

$$
\mathcal{L}_{A,r}[\psi](x) := d_A P.V. \int_{|z| \leq r} \frac{\psi(x) - \psi(x + z)}{|z|^{d+\lambda}} dz,
$$

$$
\mathcal{L}_A^s[\psi](x) := d_A \int_{|z| > r} \frac{\psi(x) - \psi(x + z)}{|z|^{d+\lambda}} dz.
$$

We shall also denote

$$
\mathcal{L}_{A,r}^s[\psi](x) := d_A \int_{r \leq |z| < s} \frac{\psi(x) - \psi(x + z)}{|z|^{d+\lambda}} dz, \text{ if } 0 < s < r; \quad \mathcal{L}_{A,r}^s[\psi](x) = 0, \text{ if } s \geq r > 0.
$$

We now recall the notion of the entropy solution for (1.1) from [9, Definition 1.2].

Definition 2.1 (Stochastic Entropy Solution). An element $u \in N^p_0(0, T, L^2(\mathbb{R}^d))$, with initial data $u_0 \in L^2(\mathbb{R}^d)$, is called a stochastic entropy solution of (1.1) if given a non-negative test function $\psi \in C^0_{\text{loc}}([0, \infty) \times \mathbb{R}^d)$ and a regular convex entropy-entropy flux pair $(\eta, \zeta)$, the following inequality holds:

$$
\int_{\mathbb{R}^d} \eta(u_0(x))\psi(0, x) \, dx + \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \left[ \eta(u(t, x))\partial_t \psi(t, x) - \nabla \psi(t, x) \cdot \zeta(u(t, x)) \right] \, dx \, dt \\
- \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \left[ \mathcal{L}_{A,r}[\psi(u(t, \cdot))](x) \psi(t, x) + A^r_\eta(u(t, x)) \mathcal{L}_{A,r}[\psi(t, \cdot)](x) \right] \, dx \, dt \\
+ \sum_{k \geq 1} \int_{\mathbb{R}^d} h_k(u(t, x))\eta'(u(t, x))\psi(t, x) \, dx \, dt + \frac{1}{2} \int_{\mathbb{R}^d} G^2(u(t, x))\eta''(u(t, x))\psi(t, x) \, dx \, dt \\
\geq 0, \text{ a.s.}
$$

2.2. Finite Difference Scheme. We begin by introducing some notations needed to define the semi-discrete finite difference scheme. Let $\Delta x$ denote a small positive number that represents the spatial discretization parameter of the numerical scheme. We set $x_\alpha = \alpha \Delta x$, for $\alpha = (\alpha_1, \ldots, \alpha_d) \in \mathbb{Z}^d$, to denote the spatial mesh points. Furthermore, following Cifani et al. [14], let us introduce the spatial grid cells

$$
R_0 = \left[-\Delta x/2, \Delta x/2\right]^d, \quad R_{\alpha} = [x_{\alpha_1-1/2}, x_{\alpha_1+1/2}) \times \cdots \times [x_{\alpha_d-1/2}, x_{\alpha_d+1/2}) = x_\alpha + R_0,
$$

where $x_{j+1/2} = x_j + \frac{\Delta x}{2}$, for $j \in \mathbb{Z}$ and $\alpha = (\alpha_1, \alpha_2, \ldots, \alpha_d) \in \mathbb{Z}^d$.

For any given $(u_\alpha) \in \mathbb{R}^{2^d}$, set $u_{\Delta x} = \sum_{\alpha \in \mathbb{Z}^d} u_\alpha \mathbbm{1}_{R_\alpha}$ (I_A denotes the characteristic function of the set A) and note that $u_{\Delta x}$ is measurable in $\mathbb{R}^d$. Moreover, $u_{\Delta x} \in L^p(\mathbb{R}^d)$ if and only if $u_{\Delta x} \in L^p(\mathbb{Z}^d)$ with

$$
p < +\infty : \quad \|u_{\Delta x}\|^p_{L^p(\mathbb{R}^d)} = \int_{\mathbb{R}^d} |u_{\Delta x}(x)|^p \, dx = \Delta x^d \sum_{\alpha \in \mathbb{Z}^d} |u_\alpha|^p = |\Delta x|^d \|(u_\alpha)|^p_{L^p(\mathbb{Z}^d)},
$$

where $\Delta x = \min_{\alpha \in \mathbb{Z}^d} \Delta x_{\alpha}$.
Furthermore, the BV semi-norm for a lattice function \( u_{\Delta x} \) is defined as

\[
|u_{\Delta x}(t, \cdot)|_{BV} = \Delta x^{d-1} \sum_{\alpha \in \mathbb{Z}^d} |u_{\alpha + e_i}(t) - u_{\alpha}(t)|,
\]

where \( \{e_1, e_2, \ldots, e_d\} \) denotes the standard basis of \( \mathbb{R}^d \). Finally, we denote by \( D_k^\pm \) the discrete forward and backward differences in space, i.e.,

\[
D_k^\pm u_{\Delta x} = \sum_{\alpha \in \mathbb{Z}^d} D_k^\pm u_{\alpha} \mathbb{I}_{R_\alpha} = \sum_{\alpha \in \mathbb{Z}^d} \pm \frac{u_{\alpha + e_k} - u_{\alpha}}{\Delta x} \mathbb{I}_{R_\alpha},
\]

where the same notation \( D_k^\pm \) is used for \( u_{\Delta x} \) and \( u_\alpha \), with a slight abuse of notation.

We now propose the following semi-discrete (in time) finite volume scheme approximating the solutions generated by the equation (1.1)

\[
du_\alpha(t) + \sum_{k=1}^{d} D_k^\mp F_k(u_\alpha, u_{\alpha + e_k}) dt + \mathbb{L}_\alpha[A(u_{\Delta x}(t, \cdot))] dt = \sigma(u_\alpha(t)) dW(t), \quad t > 0, \; \alpha \in \mathbb{Z}^d, \quad (2.2)
\]

\[
u_\alpha(0) = \frac{1}{\Delta x^d} \int_{R_\alpha} u_0(x) dx, \; \alpha \in \mathbb{Z}^d, \quad (2.3)
\]

where \( F_i \) is a monotone numerical flux corresponding to \( f_i \), for each \( 1 \leq i \leq d \). Here monotone is understood in the following sense: \( F_i \) is Lipschitz continuous from \( \mathbb{R}^2 \) to \( \mathbb{R} \), \( F_i(u, v) = f_i(u) \) for any real \( u, \; a \mapsto F_i(a, b) \) is non-decreasing and \( b \mapsto F_i(a, b) \) is non-increasing. Some classical examples of monotone flux include Engquist-Osher flux, Godunov flux, and modified Lax-Friedrichs flux. Next, we give details about the discretization of the non-local term.

### 2.2.1. Approximation of the non-local operator

Following the idea developed in [14], for any bounded \( u_{\Delta x} \), we will use the following (two layer) discretization of the non-linear non-local term: for any \( x \), \( a \mapsto F_i(a, b) \) is non-decreasing and \( b \mapsto F_i(a, b) \) is non-increasing. Some classical examples of monotone flux include Engquist-Osher flux, Godunov flux, and modified Lax-Friedrichs flux. Next, we give details about the discretization of the non-local term.

By denoting \( d\mu(z) = d\lambda \frac{dz}{|z|^{d+2 \lambda}} \) and, for any \( \alpha, \beta \in \mathbb{Z}^d \), \( G_{\alpha, \beta} = \int_{R_\alpha} \int_{|z| > \Delta x^{1/2}} [1_{R_\beta}(x) - 1_{R_\beta}(x + z)] d\mu(z) dx \), it is straightforward to verify that

\[
G_{\alpha, \beta} = G_{\alpha, \beta}, \quad G_{\alpha, \beta} \leq 0, \quad 0 \leq - \sum_{\beta \neq \alpha} G_{\alpha, \beta} = - \sum_{\beta \neq \alpha} G_{\beta, \alpha} = G_{\alpha, \alpha} \leq \Delta x^{d} \mu\left(|z| > \frac{\Delta x}{2}\right),
\]

\[
\forall k = 1, \ldots, d, \quad G_{\alpha, \beta} = G_{\alpha + e_k, \beta + e_k}.
\]

For technical reasons, we need to split the non-local integral in two parts and introduce the following notations: For any \( u_{\Delta x} = \sum_{\beta \in \mathbb{Z}^d} u_\beta \mathbb{I}_{R_\beta} \), and all \( r > \frac{\Delta x}{2} \), we write \( \mathbb{L}_\alpha[u_{\Delta x}] = \mathbb{L}_{\alpha, r}[u_{\Delta x}] + \mathbb{L}_\alpha' [u_{\Delta x}] \), where

\[
\mathbb{L}_{\alpha, r}[u_{\Delta x}] = \sum_{\alpha \in \mathbb{Z}^d} \mathbb{L}_{\alpha, r} [u(t, \cdot)]_{\alpha} \mathbb{I}_{R_\alpha} = \frac{1}{\Delta x^d} \sum_{\alpha, \beta \in \mathbb{Z}^d} G_{\alpha, \beta, r} u_\beta \mathbb{I}_{R_\alpha},
\]

with \( G_{\alpha, \beta, r} := \int_{R_\alpha} \int_{\frac{\Delta x}{2} < |z| \leq r} \left(1_{R_\beta}(x) - 1_{R_\beta}(x + z)\right) d\mu(z) dx \).
apply Itô's product rule to

In light of the above observations, we can recast the scheme (2.2) as

Remark 2.2. It is easy to check that Remark 2.1 holds for both $G_{r,\alpha}^\nu$ and $G_{\alpha,\beta,r}$, with the fact that $G_{r,\alpha}^\nu \leq \Delta x^d \mu(r > |z| > \frac{\lambda}{r^2})$ and $G_{\alpha,\beta,r} \leq \Delta x^d \mu(|z| > r)$, and that $G_{\alpha,\beta} = G_{r,\alpha}^\nu + G_{\alpha,\beta,r}$.

In light of the above observations, we can recast the scheme (2.2) as

\[
du_\alpha + \frac{1}{\Delta x} \sum_{i=1}^{d} \left[ F_i(u_{\alpha+\epsilon_i}, \alpha, \lambda, x) - F_i(u_{\alpha-\epsilon_i}, \alpha, \lambda, x) \right] dt + \frac{1}{\Delta x} \sum_{\beta \in \mathbb{Z}^d} G_{\alpha,\beta}(u_\alpha) dt = \sigma(u_\alpha) dW(t). \tag{2.5}
\]

Thanks to the assumptions on the data (A.1)–(A.4) and Remark 2.1, the solvability of (2.5) follows from a classical argument of stochastic differential equations with Lipschitz non-linearities (see e.g. Prévoit and Röckner [35, Sec. 4.1 p.55]).

2.3. Discrete entropy inequality. With the help of the above scheme (2.5), we can derive the discrete entropy inequality. To that context, let $(\eta, \zeta)$ be a $C^2$ entropy-entropy flux pair. Given a non-negative test function $\psi \in C_0^1((0, \infty) \times \mathbb{R}^d)$, consider its piecewise approximation $\psi_\alpha(t) := \frac{1}{\Delta x} \int_{R_\alpha} \psi(t, x) dx$, and apply Itô’s product rule to $\eta(u_\alpha(t)) \psi_\alpha(t)$ to yield,

\[
\eta(u_\alpha(T)) \psi_\alpha(T) - \eta(u_\alpha(0)) \psi_\alpha(0) = \int_{0}^{T} \left[ \eta(u_\alpha(t)) \partial_t \psi_\alpha(t) - \eta'(u_\alpha(t)) \left( \frac{1}{\Delta x} \sum_{i=1}^{d} \left[ F_i(u_{\alpha+\epsilon_i}, \alpha, \lambda, x) - F_i(u_{\alpha-\epsilon_i}, \alpha, \lambda, x) \right] \psi_\alpha(t) \right) \right] dt.
\]

Denote by $\overline{\psi}(t, x) = \sum_{\alpha \in \mathbb{Z}^d} \psi_\alpha(t) \mathbb{1}_{R_\alpha}(x)$. Then, integrating over $R_\alpha$ and summing over $\alpha \in \mathbb{Z}^d$ yields

\[
\int_{\mathbb{R}^d} \eta(u_\Delta x(T, x)) \overline{\psi}(T, x) dx - \int_{\mathbb{R}^d} \eta(u_\Delta x(0, x)) \overline{\psi}(0, x) dx = \int_{\Pi_T} \eta'(u_\Delta x(t, x)) \partial_t \overline{\psi}(t, x) - \sum_{\alpha \in \mathbb{Z}^d} \eta'(u_\Delta x(t, x)) \partial_t \psi_\alpha(t) dx dt
\]

\[
- \int_{\Pi_T} \eta'(u_\Delta x(t, x)) \left( \frac{1}{\Delta x} \sum_{i=1}^{d} \left[ F_i(u_{\Delta x+\epsilon_i}, x, \Delta x e_i) - F_i(u_{\Delta x-\epsilon_i}, x, \Delta x e_i) \right] \right) \overline{\psi}(t, x) dx dt
\]

\[
+ \sum_{k \geq 1} \int_{\Pi_T} g_k(u_{\Delta x}(t, x)) \eta''(u_\Delta x(t, x)) \overline{\psi}(t, x) dx dt + \frac{1}{2} \sum_{k \geq 1} \int_{\Pi_T} G^2(u_{\Delta x}(t, x)) \eta''(u_\Delta x(t, x)) \overline{\psi}(t, x) dx dt.
\]

To deal with the fractional term we follow [9, Appendix A] to notice that, for any $r > \frac{\lambda}{2}$,

\[
\int_{\Pi_T} A_{\alpha,\beta}^\nu(u_{\Delta x}(t, \cdot)) \mathbb{G}_{\lambda,\rho}^\alpha \overline{\psi}(t, \cdot) dx dt + \int_{\Pi_T} \mathbb{E}_{\alpha,\beta}^\lambda \left[ \mathbb{G}_{\lambda,\rho}^\alpha \overline{\psi}(t, \cdot) \right] dx dt.
\]
Therefore, the discrete entropy inequality is understood in the following sense:

\[ 0 \leq \int_{\mathbb{R}^d} \eta(u_{\Delta x}(0, x)) \psi(0, x) \, dx + \int_{\Pi_T} \eta(u_{\Delta x}(t, x)) \partial_t \psi(t, x) \, dx \, dt \]

\[ - \int_{\Pi_T} \eta'(u_{\Delta x}(t, x)) \frac{1}{\Delta x} \sum_{i=1}^d \left[ F_i(u_{\Delta x}(t, x), u_{\Delta x}(t, x + \Delta x e_i)) - F_i(u_{\Delta x}(t, x - \Delta x e_i), u_{\Delta x}(t, x)) \right] \psi(t, x) \, dx \, dt \]

\[ - \int_{\Pi_T} A_i^\alpha(u_{\Delta x}(t, x)) \Delta x \psi(t, x) \eta'(u_{\Delta x}(t, x)) \, dx \, dt \]

\[ + \sum_{k \geq 1} \int_{\Pi_T} g_k(u_{\Delta x}(t, x)) \eta'(u_{\Delta x}(t, x)) \psi(t, x) \, dx \, dt + \frac{1}{2} \int_{\Pi_T} \mathbb{G}^2(u_{\Delta x}(t, x)) \eta''(u_{\Delta x}(t, x)) \psi(t, x) \, dx \, dt. \]

We finish this section by stating the main result of this article, and for a proof of this main theorem we refer to Section 4.

**Theorem 2.1.** *(Main Theorem)* Let the assumptions (A.1) – (A.4) hold, and \( u_{\Delta x} \) denotes the approximate solution generated by the finite volume scheme (2.5). Moreover, let \( u \) denotes the unique BV entropy solution to the problem (1.1). Then there exists a constant \( C \), independent of \( \Delta x \), such that for all \( t \in (0, T) \)

\[ \mathbb{E} \left[ \int_{\mathbb{R}^d} \left| u_{\Delta x}(t, x) - u(t, x) \right| \, dx \right] \leq C \begin{cases} \sqrt{\Delta x}, & \text{for } \lambda < \frac{1}{2}, \\ \sqrt{\Delta x} \log \Delta x, & \text{for } \lambda = \frac{1}{2}, \\ (\Delta x)^{1-\lambda}, & \text{for } \lambda > \frac{1}{2}. \end{cases} \]

### 3. A Priori Estimates

This section is devoted to the derivation of *a priori* estimates for the approximate solutions \( u_{\Delta x}(t, x) \) under the usual assumptions.

#### 3.1. Uniform Moment Estimates

As we mentioned earlier, to ensure the convergence of the sequence of approximate solutions, one needs to obtain uniform moment estimates on it. In what follows, we start with the following simple but useful lemma which is essentially a discrete version of the entropy inequality (2.1).

**Lemma 3.1.** Let \( \eta \) be an even, \( C^2(\mathbb{R}) \) convex function with a bounded second derivative. Let \( u_\alpha(t) \) be the approximate solution generated by the finite volume scheme (2.2). Then \( u_\alpha(t) \) satisfies the following cell entropy inequality: for any \( k \in \mathbb{R} \),

\[ dq(u_\alpha(t) - k) + \frac{1}{\Delta x} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\alpha(t) - k) G_{\alpha, \beta} A(u_\beta(t)) \, dt \]

\[ + \sum_{i=1}^d \frac{|\eta'(u_\alpha(t) - k)|}{\Delta x} \left\{ \left( F_i[u_\alpha(t) \perp k, u_{\alpha+e_i}(t) \perp k] - F_i[u_{\alpha-e_i}(t) \perp k, u_\alpha(t) \perp k] \right) 
\]

\[ - \left( F_i[u_\alpha(t) \perp k, u_{\alpha+e_i}(t) \perp k] - F_i[u_{\alpha-e_i}(t) \perp k, u_\alpha(t) \perp k] \right) \right\} \, dt \]

\[ \leq \sigma(u_\alpha(t)) \eta'(u_\alpha(t) - k) \, dW(t) + \frac{1}{2} \sigma^2(u_\alpha(t)) \eta''(u_\alpha(t) - k) \, dt, \]

for all \( \alpha \in \mathbb{Z}^d \) and almost all \( \omega \in \Omega \). Here \( a \perp b := \max\{a, b\} \) and \( a \parallel b := \min\{a, b\} \).

**Proof.** A simple application of Itô’s formula applied to \( \eta(u_\alpha(t) - k) \), where \( u_\alpha(t) \) satisfies the semi-discrete finite volume scheme (2.2), leads to

\[ dq(u_\alpha(t) - k) + \frac{1}{\Delta x} \eta'(u_\alpha(t) - k) \sum_{i=1}^d \left( F_i(u_\alpha(t), u_{\alpha+e_i}(t)) - F_i(u_{\alpha-e_i}(t), u_\alpha(t)) \right) \, dt \]

\[ + \frac{1}{\Delta x} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\alpha(t) - k) G_{\alpha, \beta} A(u_\beta(t)) \, dt = \sigma(u_\alpha(t)) \eta'(u_\alpha(t) - k) \, dW(t) + \frac{1}{2} \sigma^2(u_\alpha(t)) \eta''(u_\alpha(t) - k) \, dt. \]
To manipulate terms coming from the flux $F_i$, we first observe that, for each $i = 1, 2, \ldots, d$ and any reals $u, v$ and $w$, we have

$$
\eta'(u - k) \left[ F_i(u, v) - F_i(w, u) \right] = |\eta'(u - k)| \operatorname{sign}(u - k) \left[ F_i(u, v) - F_i(w, u) \right].
$$

To simplify the notations, we denote $A := \operatorname{sign}(u - k) \left[ F_i(u, v) - F_i(w, u) \right]$. Then for $u > k$, since $F_i$ is non-decreasing with respect to its first argument and non-increasing with respect to its second one,

$$
A = F_i((u - k)^+ + k, (v - k)^+ + k) - F_i((w - k)^+ + k, (u - k)^+ + k)
$$

$$
\geq F_i((u - k)^+ + k, (v - k)^+ + k) - F_i((w - k)^+ + k, (u - k)^+ + k).
$$

Moreover,

$$
F_i[u^\perp, v^\perp] = F_i[(u - k)^- + k, (v - k)^- + k] = F_i[0, k].
$$

Therefore, we conclude that

$$
|\eta'(u - k)| A \geq |\eta'(u - k)| \left( F_i[u^\perp, v^\perp] - F_i[w^\perp, u^\perp] \right).
$$

A similar calculation reveals that the above inequality (3.2) also holds for $u \leq k$. This concludes the proof of the lemma. \hfill \Box

Now we are ready to prove uniform moment estimates. In what follows, we first state and prove the following lemma:

**Lemma 3.2.** Let the assumptions (A.1) − (A.4) hold, and $u_{\Delta x}(t, x)$ be the approximate solution generated by the semi-discrete finite volume scheme (2.2). Then, we have

$$
\sup_{\Delta x > 0} \sup_{0 \leq t \leq T} \mathbb{E} \left[ \left| u_{\Delta x}(t, \cdot) \right| \right]^{\frac{p}{p-1}} \leq M \left\| u_0 \right\|_{{L^p}(\mathbb{R}^d)}, \text{ for } p \in \mathbb{N}, \ p \geq 1.
$$

$$
\left\| u_{\Delta x}(\omega, t, \cdot) \right\|_{{L^\infty}(\mathbb{R}^d)} \leq M, \ \forall t \geq 0, \ \mathbb{P} - \text{a.s.},
$$

where $M$ is defined in assumption A.4.

**Proof.** For a given convex function as in Lemma 3.1 and $k = 0$, after taking the expectation, we have

$$
\mathbb{E}[\eta(u_0(t))] + \frac{1}{\Delta x} \int_0^t \sum_{\beta \in \mathbb{Z}^d} \mathbb{E}[\eta'(u_\alpha(s))G_{\alpha,\beta} A(u_\beta(s))] ds
$$

$$
+ \frac{1}{\Delta x} \sum_{i=1}^d \mathbb{E} \left[ \int_0^t |\eta'(u_\alpha(s))| \left( F_i[u_{\alpha + e_i}^+(s), u_{\alpha + e_i}^-(s)] - F_i[u_{\alpha - e_i}^+(s), u_{\alpha}^-(s)] - F_i[u_{\alpha - e_i}^-(s), u_{\alpha}^-] - F_i[-u_\alpha^-, u_{\alpha}^-] \right) ds \right]
$$

$$
\leq \frac{1}{2} \int_0^t \mathbb{E}[\sigma^2(u_\alpha(s))\eta''(u_\alpha(s))] ds + \eta(u_0(0)).
$$

Then, observe that the regularity of $\eta$ can be relaxed from $C^2$ to $C^1$, with $\eta'$ being Lipschitz-continuous (and $\eta'(0) = 0$). We leverage this observation to conclude

$$
\left[ \sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} |\eta'(u_\alpha)||G_{\alpha,\beta}| |A(u_\beta)| \right]^2 \leq \left[ \sum_{\beta \in \mathbb{Z}^d} |A(u_\beta)| \right] \left[ \sum_{\alpha \in \mathbb{Z}^d} |\eta'(u_\alpha)||G_{\alpha,\beta}| \right]^2
$$

$$
\leq C(A, \eta)||u_\alpha(0)||_{L^2(\mathbb{R}^d)}^2 \sum_{\alpha \in \mathbb{Z}^d} |u_\alpha(0)|^2 \sum_{\beta \in \mathbb{Z}^d} \left[ \sum_{\alpha \in \mathbb{Z}^d} |G_{\alpha,\beta}|^2 \right] \leq C(A, \eta)||u_\alpha(0)||_{L^2(\mathbb{R}^d)}^2 \sum_{\alpha \in \mathbb{Z}^d} |u_\alpha(0)|^2 \sum_{\beta \in \mathbb{Z}^d} \left[ \sum_{\alpha \in \mathbb{Z}^d} |G_{\alpha,\beta}|^2 \right] < +\infty, \ \mathbb{P} - \text{a.s.}
$$

(3.6)
Moreover, making use of Remark 2.1 and (3.6), one has that
\[
\sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\alpha) G_{\alpha, \beta} A(u_\beta) = \sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\alpha) G_{\beta, \alpha} (A(u_\beta) - A(u_\alpha))
\]
\[
\sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\beta) G_{\alpha, \beta} A(u_\alpha) = \sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\beta) G_{\alpha, \beta} (A(u_\alpha) - A(u_\beta)) = - \sum_{\alpha, \beta \in \mathbb{Z}^d} \eta'(u_\beta) G_{\alpha, \beta} (A(u_\beta) - A(u_\alpha)).
\]
This implies that
\[
\sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} \eta'(u_\alpha) G_{\alpha, \beta} A(u_\beta) = \frac{1}{2} \sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} (\eta'(u_\alpha) - \eta'(u_\beta)) G_{\alpha, \beta} (A(u_\beta) - A(u_\alpha)) \geq 0. \quad (3.7)
\]
Assuming that \(\eta(0) = 0\) and since \((u_\alpha) \in C([0, T], \ell^2(\mathbb{Z}^d))\), one gets
\[
\sum_{\alpha \in \mathbb{Z}^d} E[\eta(u_\alpha(t))] + \int_0^T \sum_{\alpha \in \mathbb{Z}^d} \sum_{\beta \in \mathbb{Z}^d} E[\eta'(u_\alpha(s)) G_{\alpha, \beta} A(u_\beta(s))] ds \leq \frac{1}{2} \sum_{i=1}^d \sum_{\alpha \in \mathbb{Z}^d} \int_0^T |\eta'(u_\alpha(s))| \left( \left| F_i[u_\alpha^+(s), u_{\alpha+e_i}^+(s)] - F_i[u_{\alpha-e_i}^-(s), u_{\alpha}^-(s)] \right| - \left| F_i[-u_{\alpha}^-(s), -u_{\alpha-e_i}^+(s)] - F_i[-u_{\alpha-e_i}^-(s), -u_{\alpha}^-(s)] \right| \right) ds + \sum_{\alpha \in \mathbb{Z}^d} \eta(u_\alpha(0)). \quad (3.8)
\]
Consider, in (3.8), that \(\eta = \eta_\delta\) is the convex even function such that \(\eta_\delta(0) = 0\) and \(\eta_\delta(x) = \min(1, \frac{x}{\delta})\) for positive \(x\). Noting that
\[
\left( F_i[0, u_{\alpha+e_i}^+(s)] - F_i[u_{\alpha-e_i}^-(s), 0] \right) - \left( F_i[0, -u_{\alpha+e_i}^-(s)] - F_i[-u_{\alpha-e_i}^-(s), 0] \right) \leq 0,
\]
passing to the limit \(\delta \to 0\), one gets for any \(t\),
\[
0 \geq \sum_{\alpha} E[u_\alpha(t)] - \sum_{\alpha} |u_\alpha(0)|
\]
\[
+ \frac{1}{\Delta x} \sum_{i=1}^d \sum_{\alpha} \int_0^T \left\{ \left( F_i[u_{\alpha}^+(s), u_{\alpha+e_i}^+(s)] - F_i[u_{\alpha-e_i}^-(s), u_{\alpha}^-(s)] \right) - \left( F_i[-u_{\alpha}^-(s), -u_{\alpha+e_i}^+(s)] - F_i[-u_{\alpha-e_i}^-(s), -u_{\alpha}^-(s)] \right) \right\} ds.
\]
Multiplying the above inequality by \(\Delta x^d\), we are left with
\[
\Delta x^d \sum_{\alpha \in \mathbb{Z}^d} E[u_\alpha(t)] \leq \Delta x^d \sum_{\alpha \in \mathbb{Z}^d} |u_\alpha(0)| \text{ and } \sup_{\Delta x > 0} \sup_{0 \leq t \leq T} E\left[ \|u_{\Delta x}(t, \cdot)\|_{L^1(\mathbb{R}^d)} \right] \leq \|u_0\|_{L^1(\mathbb{R}^d)}.
\]
In order to prove the maximum principle, let us assume for the moment that, for any \(i \in \{1, \ldots, d\}\) and any real \(u\), \(A(u)\) is replaced by \(A(T_M(u))\) and \(f_i(u)\) by \(f_i(T_M(u))\) where \(T_M\) denotes the classical truncation: \(T_M(u) = \max(0, \min(u, M))\); and that the numerical flux becomes \(F_i(T_M(u), T_M(u))\). This is a monotone numerical flux associated with \(f_i(T_M(u))\) and everything that has been done so far remains valid.

Note that as soon as the \(L^\infty\) estimate in (3.4) will be proved with the perturbation \(A(T_M(u))\), \(f_i(T_M(u))\) and \(F_i(T_M(u), T_M(u))\), then the corresponding solution will be a solution to (2.2), and, as the latter is unique, the estimate will hold for the solution \((u_\alpha)\).

Under the conditions described above, apply Itô formula to \(\eta(u_\alpha(t) - M)\), where \(\eta((-\infty, 0]) = \{0\}\) and \(u_\alpha(t)\) satisfies the semi-discrete finite volume scheme (2.2).

As \(\eta\) is convex, \(\eta' \geq 0\), one gets that
\[
\sigma(u_\alpha(t)) \eta'(u_\alpha(t) - M) = 0, \quad \sigma^2(u_\alpha(t)) \eta''(u_\alpha(t) - M) = 0
\]
and
Proof. Let us write 
\[ \eta(t) = \left( F_i(u_\alpha(t), u_{\alpha+e_i}(t)) - F_i(u_{\alpha-e_i}(t), u_\alpha(t)) \right) \]
so that, for a suitable regular approximation of \( \eta \) and one is interested, via Itô's formula, to estimate \( \eta(u_{\alpha+e_j} - u_\alpha) \) and more precisely, the limit when \( \eta \) converges to the absolute value function.

Note, by Remark 2.1, that
\[ \sum_{\beta \in \mathbb{Z}^d} G_{\alpha,\beta} A(u_{\beta}) = \sum_{\beta \in \mathbb{Z}^d} G_{\alpha,\beta} A(u_{\beta}) - \sum_{\beta \in \mathbb{Z}^d} G_{\alpha,\beta} A(u_{\beta}) = \sum_{\beta \in \mathbb{Z}^d} G_{\alpha,\beta} [A(u_{\beta + e_j}) - A(u_{\beta})], \]
so that, for a suitable regular approximation of \( \eta = \eta_\delta \), one has, for any \( t \), that
\[ E\eta(u_{\alpha+e_j} - u_\alpha(t)) = \frac{1}{\Delta x} \sum_{\beta \in \mathbb{Z}^d} E \int_0^t \eta'(u_{\alpha+e_j} - u_\alpha)(s) G_{\alpha,\beta} [A(u_{\beta + e_j}) - A(u_{\beta})](s) ds \]
\[ = -\sum_{i=1}^d E \int_0^t \eta'(u_{\alpha+e_j} - u_\alpha)(s) [F_i(u_{\alpha+e_j}, u_{\alpha+e_i+e_j}) - F_i(u_{\alpha-e_i+e_j}, u_{\alpha+e_j}) - F_i(u_{\alpha+e_i}, u_{\alpha+e_j}) + F_i(u_{\alpha-e_i}, u_{\alpha})](s) ds \]
\[ + \frac{1}{2} E \int_0^t \eta''(u_{\alpha+e_j} - u_\alpha)(s) \sigma(u_{\alpha+e_j}) - \sigma(u_{\alpha})(s) dy + E\eta(u_{\alpha+e_j} - u_\alpha)(0). \]
Using Lebesgue's theorem,
\[ \lim_{\delta \to 0} E \int_0^t \eta''(u_{\alpha+e_j} - u_\alpha)(s) \sigma(u_{\alpha+e_j}) - \sigma(u_{\alpha})(s) dy = 0, \]
and one is able to pass to the limit on \( \eta_\delta \) and replace, above, \( \eta(u) \) by \( |u| \) and \( \eta'(u) \) by \( \text{sgn}(u) \).

Note that, by adding and subtracting appropriate terms, we have
\[ F_i(u_{\alpha+e_j}, u_{\alpha+e_i+e_j}) - F_i(u_{\alpha-e_i+e_j}, u_{\alpha+e_j}) - F_i(u_{\alpha+e_i}, u_{\alpha+e_j}) + F_i(u_{\alpha-e_i}, u_{\alpha}) \]
}\[ \]
\[ \geq 0. \]

By argument similar to (3.7), \( \eta(u_\alpha(t) - M) \leq \eta(u_\alpha(0) - M) = 0 \) and u_\alpha \leq M. Moreover, applying Itô formula to \( \eta(u_\alpha(t) + M) \) implies that u_\alpha \geq -M. Finally, assuming \( p > 1 \), the result is proved by using an argument of interpolation.

3.2. Spatial Bounded Variation. Like its deterministic counterpart, we derive spatial BV bound for the approximate solutions under usual assumptions.

Lemma 3.3. Let the assumptions be true. Let \( u_{\Delta x}(t, x) \) be the finite volume approximations prescribed by the finite difference scheme (2.2). Then for any \( t > 0 \)
\[ E\left[ |u_{\Delta x}(t, \cdot)|_{BV(\mathbb{R}^d)} \right] \leq C E\left[ |u_0(\cdot)|_{BV(\mathbb{R}^d)} \right]. \] (3.9)

Proof. One has
\[ du_\alpha + \frac{1}{\Delta x} \sum_{i=1}^d \left[ F_i(u_\alpha, u_{\alpha+e_i}) - F_i(u_{\alpha-e_i}, u_\alpha) \right] dt + \frac{1}{\Delta x^d} \sum_{\beta \in \mathbb{Z}^d} G_{\alpha,\beta} A(u_{\beta}) dt = \sigma(u_\alpha) dW(t), \]
\[ du_{\alpha+e_j} + \frac{1}{\Delta x} \sum_{i=1}^d \left[ F_i(u_{\alpha+e_j}, u_{\alpha+e_i+e_j}) - F_i(u_{\alpha-e_i+e_j}, u_{\alpha+e_j}) \right] dt + \frac{1}{\Delta x^d} \sum_{\beta \in \mathbb{Z}^d} G_{\alpha+e_j,\beta} A(u_{\beta}) dt = \sigma(u_{\alpha+e_j}) dW(t), \]
and one is interested, via Itô’s formula, to estimate \( \eta(u_{\alpha+e_j} - u_\alpha) \) and more precisely, the limit when \( \eta \) converges to the absolute value function.

1 the convex even function such that \( \eta_\delta(0) = 0 \) and \( \eta_\delta'(x) = \min(1, \frac{x}{\delta}) \) for positive \( x \), \( 0 \leq \eta(u) \leq \frac{1}{\delta} u^2 \) and \( u^2 \eta_\delta''(u) \leq \eta_\delta(u) \).
Therefore, using that $F_i$ is non-decreasing in its first argument and non-increasing in its second one,

\[ E \{ - \text{sgn}(u_{a+e_j} - u_a) \left[ F_i(u_{a+e_j}, u_{a+e_j}) - F_i(u_{a-e_j}, u_a) - F_i(u_{a-e_j}, u_{a+e_j}) + F_i(u_{a-e_j}, u_a) \right] \} \]

\[ \leq E \{ - |\eta'(u_{a+e_j} - u_a)||F_i(u_{a+e_j}, u_{a+e_j}) - F_i(u_{a, u_{a+e_j}})|u_{a+e_j} - u_a| \}
\]

Hence, after summing over $j$, we get

\[ E \{ \sum_{j=1}^{d} \left[ u_{a+e_j}(t) - u_a(t) \right] \} \leq E \{ \sum_{j=1}^{d} \left[ u_{a+e_j}(0) - u_a(0) \right] \} . \]

Note that, in view of the lower semi-continuity property and the positivity of the total variation $TV_x$, $E[TV_x(u)]$ makes sense for any $u \in L^1(\Omega \times \mathbb{R})$. Since $u_0 \in BV(\mathbb{R})$, we conclude that, for all $t > 0$

\[ E \{ TV_x(u\Delta_t(t)) \} \leq E \{ TV_x(u_0) \} . \]
Again, since $E\left[\|u_{\Delta}(t,\cdot)\|_{L^1(\mathbb{R})}\right] \leq C E\left[\|u_0(\cdot)\|_{L^1(\mathbb{R}^d)}\right]$, we arrive at the following conclusion that the approximate solution $u_{\Delta}(t,x)$ lies in the spatial BV class and satisfies (3.9). This completes the proof.

\[ \square \]

4. Proof of the Main Result: Theorem 2.1

We start by introducing a special class of entropy functions, called convex approximation of absolute value function. To do so, let $\eta : \mathbb{R} \to \mathbb{R}$ be a $C^\infty$ function satisfying

$$
\eta(0) = 0, \quad \eta(-r) = \eta(r), \quad \eta'(r) = -\eta'(r), \quad \eta'' \geq 0,
$$

and

$$
\eta'(r) = \begin{cases} 
-1, & \text{when } r \leq -1, \\
\in [-1,1], & \text{when } |r| < 1, \\
+1, & \text{when } r \geq 1.
\end{cases}
$$

For any $\varepsilon > 0$, define $\eta_{\varepsilon} : \mathbb{R} \to \mathbb{R}$ by $\eta_{\varepsilon}(r) = \varepsilon \eta(\frac{r}{\varepsilon})$. Then

$$
|\eta'| - L_1 \varepsilon \leq \eta_{\varepsilon}(r) \leq |\eta'| \quad \text{and} \quad |\eta_{\varepsilon}''(r)| \leq \frac{L_2}{\varepsilon} 1_{|r| \leq \varepsilon},
$$

where $L_1 := \sup_{|r| \leq 1} |\eta' - \eta(r)|$ and $L_2 := \sup_{|r| \leq 1} |\eta''(r)|$.

Moreover, for $\eta = \eta_{\varepsilon}$, we define

$$
Q_k^a(a,b) = \int_0^a \eta'(r) f_k^a(r) dr,
Q_k^b(a) = (f_k^a(a,b), f_k^2(a,b), \ldots, f_k^n(a,b)),
Q(a,b) = \text{sign}(a-b)(f(a) - f(b)) = f(a+b) - f(b).
$$

For a small positive number $\varepsilon > 0$, we consider the following parabolic perturbation of (1.1)

$$
d u_{\varepsilon}(t,x) - \varepsilon \Delta u_{\varepsilon}(t,x) dt + \mathcal{L}[A(u_{\varepsilon}(t,\cdot))](x) dt - \text{div}_x f(u_{\varepsilon}(t,x)) dt = \sigma(u_{\varepsilon}(t,x)) dW(t).
$$

Following [9], we remark that it has a unique weak solution $u_{\varepsilon}(t,x)$ with initial data $u_{\varepsilon}(0,x) = u_0^\varepsilon(x) \in H^1(\mathbb{R}^d)$, where $u_0^\varepsilon$ converges to $u_0$ in $L^2(\mathbb{R}^d)$. Notice that $u_{\varepsilon} \in H^1(\mathbb{R}^d)$, while for technical reasons we require higher regularity of $u_{\varepsilon}$, therefore, we need to regularize $u_{\varepsilon}$ by a space convolution. Let $\{\rho_\gamma\}_\gamma$ be a given mollifier-sequence in $\mathbb{R}^d$. Then following [9] we observe that $u_{\varepsilon}^\gamma := u_{\varepsilon} \ast \rho_\gamma$ satisfies

$$
\partial_t [u_{\varepsilon} \ast \rho_\gamma - \int_0^t \sigma(u_{\varepsilon} \ast \rho_\gamma) dW] - [\varepsilon \Delta (u_{\varepsilon} \ast \rho_\gamma) - \mathcal{L}[A(u_{\varepsilon} \ast \rho_\gamma)] + \text{div}(f(u_{\varepsilon} \ast \rho_\gamma))] = 0.
$$

Let $\sigma$ be the standard nonnegative mollifiers on $\mathbb{R}$ and $\mathbb{R}^d$ respectively such that $\text{supp}(\rho) \subset [-1,0]$ and $\text{supp}(\varphi) = \overline{B}_1(0)$. We define $\rho_{\delta_0}(r) = \frac{1}{2\delta_0} \rho\left(\frac{r}{\delta_0}\right)$ and $g_\delta(x) = \frac{1}{\delta} g\left(\frac{x}{\delta}\right)$, where $\delta$ and $\delta_0$ are two positive parameters. Given a nonnegative test function $\psi \in C^2(0,\infty) \times \mathbb{R}^d$ and two positive constants $\delta$ and $\delta_0$, we define

$$
\varphi_{\delta,\delta_0}(t,x,s,y) = \rho_{\delta_0}(t-s) g_{\delta}(x-y) \psi(t,x).
$$

Clearly $\rho_{\delta_0}(t-s) \neq 0$ only if $s - \delta_0 \leq t \leq s$ and hence $\varphi_{\delta,\delta_0}(t,x,s,y) = 0$, outside $s - \delta_0 \leq t \leq s$. Moreover, let $\varsigma$ be the standard symmetric nonnegative mollifier on $\mathbb{R}$ with support in $[-1,1]$ and $\varsigma(r) = \frac{1}{2} \varsigma\left(\frac{r}{\delta}\right)$, for $l > 0$.

We multiply the entropy inequality (2.6) by $\varsigma(u_{\Delta y}(s,y) - k)$, take the expectation and integrate with respect to $s$ and $k$ over $\Pi_T \subset \mathbb{R} \times \mathbb{R}$ to get the following form,

$$
0 \leq E \left[ \int_{\mathbb{R} \times \Pi_T \times \mathbb{R}^d} \eta(u_{\varepsilon}^\gamma(0,x) - k) \varphi_{\delta,\delta_0}(0,x,s,y) \varsigma(u_{\Delta y}(s,y) - k) dx dy ds dk \right] 
+ E \left[ \int_{\mathbb{R} \times \Pi_T \times \Pi_T} \eta(u_{\varepsilon}^\gamma(t,x) - k) \partial_t \varphi_{\delta,\delta_0}(t,x,s,y) \varsigma(u_{\Delta y}(s,y) - k) dt dy ds dk \right] 
+ E \left[ \sum_{n \geq 1} \int_{\Pi_T} g_n(u_{\varepsilon}^\gamma(t,x)) \eta'(u_{\varepsilon}^\gamma(t,x) - k) \varphi_{\delta,\delta_0}(t,x,s,y) d\beta_k(t) dx \right] \varsigma(u_{\Delta y}(s,y) - k) dy ds dk.
$$
\[ E \left[ \int_{\mathbb{R} \times \mathbb{R}^d} g_\delta(u_\gamma(t, x, y)) \varphi(y) \, dy \int_{\mathbb{R}^d} \eta(y) \right] \leq \frac{1}{2} \mathbb{E} \left[ \int_{\mathbb{R} \times \mathbb{R}^d} \mathbb{G}^2(u_\gamma(t, x))^2 \, dx \, dy \, dt \, dk \right] + \mathbb{E} \left[ \int_{\mathbb{R} \times \mathbb{R}^d} \mathbb{G}^2(u_\gamma(t, x))^2 \, dx \, dy \, dt \, dk \right] - \mathbb{E} \left[ \int_{\mathbb{R} \times \mathbb{R}^d} \mathbb{G}^2(u_\gamma(t, x))^2 \, dx \, dy \, dt \, dk \right] - \sum_{j=1}^{\infty} \mathbb{E} \left[ \int_{\mathbb{R} \times \mathbb{R}^d} \mathbb{G}^2(u_\gamma(t, x))^2 \, dx \, dy \, dt \, dk \right] := I_1 + I_2 + I_3 + I_4 + I_5 + I_6 + I_7 + I_8. \]
Finally, it follows that
\[
\lim_{t \to 0} \lim_{\eta \to 0} \lim_{\gamma \to 0, \delta_0 \to 0} (I_3 + J_3 + I_4 + J_4) = 0.
\]

Next, we move on to estimate the terms coming from the associated flux function.

**Lemma 4.3.** It holds that
\[
\limsup_{t \to 0} \lim_{\eta \to 0} \lim_{\gamma \to 0, \delta_0 \to 0} (I_5 + J_5) \leq E \left[ \int_{R^d} \int_{R^d} Q(u_{\epsilon}(t, x), u_{\Delta y}(t, y)) \cdot \nabla \psi(t, x) g_\delta(x - y) \; dy \; dx \; dt \right] + \frac{\Delta y}{\delta} C\|u_0\|_{BV} \int_0^T \|\psi(t)\|_\infty dt.
\]

**Proof.** We start with
\[
J_5 = -E \left[ \int_{R^d} \int_{R^d} \int_{R^d} \eta_1(u_{\Delta y}(s, y) - k) \frac{1}{\Delta y} \sum_{i=1}^d \left[ F_i(u_{\Delta y}(s, y)) - F_i(u_{\Delta y}(s, y + \Delta y e_i)) \right] \phi_{\delta, \gamma}(s, y) \Psi(t, y) \; ds \; dy \; dt \right]
\]
\[
\delta_0 \to 0 - E \left[ \int_{R^d} \int_{R^d} \int_{R^d} \eta_1(u_{\Delta y}(t, y) - k) \frac{1}{\Delta y} \sum_{i=1}^d \left[ F_i(u_{\Delta y}(t, y)) - F_i(u_{\Delta y}(t, y + \Delta y e_i)) \right] \rho_{\delta}(x - y) \Psi(t, y) \; dx \; dy \; dt \right]
\]
\[
\gamma \to 0 - E \left[ \int_{R^d} \int_{R^d} \int_{R^d} \eta_1(u_{\Delta y}(t, y) - k) \frac{1}{\Delta y} \sum_{i=1}^d \left[ F_i(u_{\Delta y}(t, y)) - F_i(u_{\Delta y}(t, y + \Delta y e_i)) \right] \rho_{\delta}(x - y) \Psi(t, y) \; dx \; dy \; dt \right]
\]
\[
\eta \to 0 - E \left[ \int_{R^d} \int_{R^d} \int_{R^d} \eta_1(u_{\Delta y}(t, y) - k) \frac{1}{\Delta y} \sum_{i=1}^d \left[ F_i(u_{\Delta y}(t, y)) - F_i(u_{\Delta y}(t, y + \Delta y e_i)) \right] \rho_{\delta}(x - y) \Psi(t, y) \; dx \; dy \; dt \right].
\]

By arguments similar to the proof of Lemma 3.1 and since, for any \(v, w\),
\[
F_i(k, k \wedge v) - F_i(k \wedge w, k) - \left[ F_i(k, k \perp v) - F_i(k \perp w, k) \right] \leq 0,
\]
one gets that
\[
- \sum_{i=1}^d \left[ F_i(u_{\Delta y}(t, y)) - F_i(u_{\Delta y}(t, y + \Delta y e_i)) \right] \rho_{\delta}(x - y) \; dy \leq \int_{R^d} \sum_{i=1}^d \left[ F_i(u_{\Delta y}(t, y)) - F_i(u_{\Delta y}(t, y + \Delta y e_i)) \right] \rho_{\delta}(x - y) \; dy.
\]

Thus,
\[
\int_{R^d} \sum_{i=1}^d \left[ F_i(u_{\Delta y}(t, y)) - F_i(u_{\Delta y}(t, y + \Delta y e_i)) \right] \rho_{\delta}(x - y) \; dy \leq \int_{R^d} \sum_{i=1}^d \left[ F_i(k \wedge u_{\Delta y}(t, y)) - F_i(k \perp u_{\Delta y}(t, y)) \right] \rho_{\delta}(x - y) \; dy.
\]
\[-\int_{\mathbb{R}^d} \left[ f_i(k \mathcal{T} u_{\Delta y}(t, y)) - f_i(k \mathcal{T} u_{\Delta y}(t, y - \Delta y e_i)) \right] \rho_\delta(x - y) \, dy \]
\[+ \sum_{a \in \mathbb{Z}^d} \left[ f_i(k \mathcal{T} u_{\alpha}(t), k \perp u_{\alpha + e_i}(t)) - f_i(k \mathcal{T} u_{\alpha}(t)) \right] \int_{\mathcal{R}_a} \rho_\delta(x - y) \, dy - \int_{\mathcal{R}_{\alpha + e_i}} \rho_\delta(x - y) \, dy \]
\[+ \sum_{a \in \mathbb{Z}^d} \left[ f_i(k \mathcal{T} u_{\alpha}(t), k \mathcal{T} u_{\alpha + e_i}(t)) - f_i(k \mathcal{T} u_{\alpha}(t)) \right] \int_{\mathcal{R}_a} \rho_\delta(x - y) \, dy - \int_{\mathcal{R}_{\alpha + e_i}} \rho_\delta(x - y) \, dy \]
\[\leq \int_{\mathbb{R}^d} \left[ f_i(k \mathcal{T} u_{\Delta y}(t, y)) - f_i(k \mathcal{T} u_{\Delta y}(t, y - \Delta y e_i)) \right] \rho_\delta(x - y) \, dy \]
\[+ \sum_{a \in \mathbb{Z}^d} (F_i(u_{\alpha + e_i}(t)) - u_{\alpha}(t)) \int_{\mathcal{R}_a} |\rho_\delta(x - y) - \rho_\delta(x - y - \Delta y e_i)| \, dy. \]

Since

\[E \left[ \int_{\mathbb{R}^d} \sum_{a \in \mathbb{Z}^d} |u_{\alpha + e_i}(t) - u_{\alpha}(t)| \int_{\mathcal{R}_a} |\rho_\delta(x - y) - \rho_\delta(x - y - \Delta y e_i)| \, dy \right] \leq 1 \Delta y \int_{\mathbb{R}^d} \left[ \int_{\mathcal{R}_a} \rho_\delta(x - y) \, dy \right] \psi(t, x) \, dx \, dt \]

we are left to

\[\lim_{\delta \to 0} J_5 \leq C \|u_0\|_{BV} \int_0^T \|\psi(t)\|_{\infty} dt \frac{\Delta y}{\delta} \]
\[+ \frac{1}{\Delta y} E \left[ \int_{\mathbb{R}^d} \left[ f_i(k \mathcal{T} u_{\Delta y}(t, y)) - f_i(k \mathcal{T} u_{\Delta y}(t, y - \Delta y e_i)) \right] \rho_\delta(x - y) \, dy \right] \times \psi(t, x) \mathcal{Q}(u_{\epsilon}(t, x) - k) \, dx \, dt \, dk \]
\[+ \frac{1}{\Delta y} E \left[ \int_{\mathbb{R}^d} \left[ f_i(k \mathcal{T} u_{\Delta y}(t, y)) - f_i(k \mathcal{T} u_{\Delta y}(t, y - \Delta y e_i)) \right] \rho_\delta(x - y) \, dy \right] \times \psi(t, x) \mathcal{Q}(u_{\epsilon}(t, x) - k) \, dx \, dt \, dk \]
\[= \frac{C \Delta y \|u_0\|_{BV}}{\delta} \int_0^T \|\psi(t)\|_{\infty} dt \]
\[+ \frac{1}{\Delta y} E \left[ \int_{\mathbb{R}^d} \left[ \{Q(u_{\Delta y}(t, y), k)\} - \{Q(u_{\Delta y}(t, y - \Delta y e_i), k)\} \right] \rho_\delta(x - y) \, dy \right] \times \psi(t, x) \mathcal{Q}(u_{\epsilon}(t, x) - k) \, dx \, dt \, dk \].

Here \(\{Q(u, v)\}\) denotes the \(i\)-th component of \(Q(u, v)\). Next, passing to the limit over \(t\) yields

\[\limsup_{t \to \infty} \lim_{\delta \to 0} J_5 \leq \frac{\Delta y C \|u_0\|_{BV}}{\delta} \int_0^T \|\psi(t)\|_{\infty} dt \]
\[- \frac{1}{\Delta y} E \left[ \int_{\mathbb{R}^d} \left[ \{Q(u_{\Delta y}(t, y), u_{\epsilon}(t, x))\} - \{Q(u_{\Delta y}(t, y - \Delta y e_i), u_{\epsilon}(t, x))\} \right] \rho_\delta(x - y) \, dy \psi(t, x) \, dx \right]. \]
Note that,
\[- \frac{1}{\Delta y} \int_{\mathbb{R}^d} \left\{ Q(u_{\Delta y}(t, y), u_z(t, x)) \right\}_i - \left\{ Q(u_{\Delta y}(t, y - \Delta y e_i), u_z(t, x)) \right\}_i \rho_5(x - y) \, dy\]
\[= - \frac{1}{\Delta y} \int_{\mathbb{R}^d} \left[ Q(u_{\Delta y}(t, y), u_z(t, x)) \right] \left[ \rho_5(x - y) - \rho_5(x - y - \Delta y e_i) \right] \, dy\]
\[= - \int_{\mathbb{R}^d} \left\{ Q(u_{\Delta y}(t, y), u_z(t, x)) \right\}_i \left[ \frac{\rho_5(x - y) - \rho_5(x - y - \Delta y e_i)}{\Delta y} \right] - \partial_t \rho_5(x - y) + \partial_i \rho_5(x - y) \, dy.\]

Then,
\[
\int_{\mathbb{R}^d} \left\{ Q(u_{\Delta y}(t, y), u_z(t, x)) \right\}_i \left[ \frac{\rho_5(x - y) - \rho_5(x - y - \Delta y e_i)}{\Delta y} \right] - \partial_t \rho_5(x - y) \, dy
\]
\[= \int_{\mathbb{R}^d} \left[ \frac{\rho_5(x - y) - \rho_5(x - y - \Delta y e_i)}{\Delta y} \right] - \partial_t \rho_5(x - y) \, dy
\]
\[= \int_{\mathbb{R}^d} \int_0^0 \left[ Q(u_{\Delta y}(t, y + \tau e_i), u_z(t, x)) \right]_i - Q(u_{\Delta y}(t, y), u_z(t, x)) \right\}_i \partial_t \rho_5(x - y) \, d\tau \, dy
\]
\[\leq \| f \|_{L_p} \int_{\mathbb{R}^d} \int_0^0 |u_{\Delta y}(t, y + \tau e_i) - u_{\Delta y}(t, y)| \left| \partial_t \rho_5(x - y) \right| d\tau \, dy.\]

So that
\[
\frac{\| f \|_{L_p}}{\Delta y} \int_{\mathbb{R}^d} \sum_{i=1}^d \int_0^T \int_{\mathbb{R}^d} \left| u_{\Delta y}(t, y + \tau e_i) - u_{\Delta y}(t, y) \right| \left| \partial_t \rho_5(x - y) \right| d\tau \, dy \, dt
\]
\[\leq \frac{\| f \|_{L_p}}{\Delta y} \sum_{i=1}^d \int_0^T \left\| \psi(t) \right\|_\infty \int_{\mathbb{R}^d} \int_0^0 \left| u_{\Delta y}(t, y + \tau e_i) - u_{\Delta y}(t, y) \right| d\tau \, dy \, dt
\]
\[\leq \frac{\| f \|_{L_p}}{\Delta y} \sum_{i=1}^d \int_0^T \left\| \psi(t) \right\|_\infty \Delta y^d \sum_{i=1}^d \left| u_{\alpha - e_i}(t) - u_{\alpha}(t) \right| d\tau \, dt
\]
\[\leq \frac{\| f \|_{L_p} \Delta y}{\delta} \| u_0 \|_{BV} \int_0^T \left\| \psi(t) \right\|_\infty \, dt.\]

In conclusion,
\[
\lim_{\delta \to 0} \lim_{\delta \to 0} \lim_{\eta \to 0} J_5
\]
\[\leq \frac{\Delta y}{\delta} C\| u_0 \|_{BV} \int_0^T \left\| \psi(t) \right\|_\infty \, dt - \frac{1}{\Delta y} \int_{\Pi_T \times \mathbb{R}^d} Q(u_{\Delta y}(t, y), u_z(t, x)) \cdot \nabla \rho_5(x - y) \psi(t, x) \, dy \, dx \, dt.
\]

For the other term, we simply follow [2] to conclude,
\[
\lim_{\delta \to 0} \lim_{\delta \to 0} \lim_{\eta \to 0} I_5 = \mathbb{E} \left[ \int_{\Pi_T} \int_{\mathbb{R}^d} Q(u_z(t, x), u_{\Delta y}(t, y)) \cdot \nabla_x (\rho_5(x - y) \psi(t, x)) \, dy \, dx \, dt \right],
\]
and the lemma is proved.

\begin{lemma}
[2] It holds that
\[
\lim_{\varepsilon \to 0} \lim_{l \to 0} \lim_{\eta \to 0} \lim_{\delta \to 0} \sup |I_8| = 0.
\]
\end{lemma}

We now estimate the terms coming from the fractional operator in the next two lemmas. We choose \( r = \Delta y \) for the subsequent calculations.
Lemma 4.5. It holds that

\[
\lim_{t \to 0} \sup_{\eta \to \cdot} \lim_{\delta \to 0} \sup(I_0 + J_0) \leq \mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} |A(u_{\Delta y}(t, y)) - A(u_{\varepsilon}(t, x))| \mathcal{L}_\lambda^v[\psi(t, \cdot)](x) \rho_s(x - y) \, dy \, dx \, dt \right] \\
+ \frac{C_3}{\delta} \|A'\|_{\infty} \|\psi\|_{\infty} \|u_0\|_{B^V} \begin{cases} \\
\Delta y, & \text{if } \lambda < \frac{1}{2}, \\
|\Delta y| \ln |\Delta y|, & \text{if } \lambda = \frac{1}{2}, \\
|\Delta y|^{2(1 - \lambda)}, & \text{if } \lambda > \frac{1}{2}.
\end{cases}
\]

Proof. We have

\[
I_0 + J_0 = - \mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} \mathcal{L}_\lambda^v[A(u_{\Delta y}(s, \cdot))](y) \, dy \, dx \, dt \right] \\
- \mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} \mathcal{L}_\lambda^c[A(u_{\varepsilon}(t, \cdot))](x) \phi_{\delta, \delta_0}(t, x, s, y) \eta'(u_{\Delta y}(s, y) - k) \, dx \, dt \, ds \right] \\
- \mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} \mathcal{L}_\lambda^c[A(u_{\varepsilon}(t, \cdot))](x) \phi_{\delta, \delta_0}(t, x, s, y) \eta'(u_{\Delta y}(s, y) - k) \, dx \, dt \, ds \right] \\
- \mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} \mathcal{L}_\lambda^c[A(u_{\Delta y}(s, \cdot))](y) \phi_{\delta, \delta_0}(t, x, s, y) \eta'(u_{\Delta y}(s, y) - k) \, dy \, dx \, dt \right] \\
\leq A + B + C.
\]

Following [9, Lemma 3.4, Lemma 4.6] we have,

\[
\lim_{\eta \to \cdot} \lim_{\delta \to 0} (A + B) \leq -\mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} |A(u_{\Delta y}(t, y)) - A(u_{\varepsilon}(t, x))| \mathcal{L}_\lambda^v[\psi(t, \cdot)](x) \rho_s(x - y) \, dx \, dt \, dy \right] \\
+ C(\lambda) \frac{\|A'\|_{\infty}}{\rho^{2\lambda}}.
\]

Thus we conclude that

\[
\lim_{t \to 0} \sup_{\eta \to \cdot} \lim_{\delta \to 0} (A + B) \leq -\mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} |A(u_{\Delta y}(t, y)) - A(u_{\varepsilon}(t, x))| \mathcal{L}_\lambda^v[\psi(t, \cdot)](x) \rho_s(x - y) \, dx \, dt \, dy \right].
\]

We now consider,

\[
C = -\mathbb{E} \left[ \int_{\mathbb{R}^d \times \Pi_T} \mathcal{L}_\lambda^v[A(u_{\Delta y}(s, \cdot))](y) \phi_{\delta, \delta_0}(t, x, s, y) \eta'(u_{\Delta y}(s, y) - k) \, dy \, dx \, dt \right] \\
= \mathbb{E} \left[ \int_{\Pi_T \times \Pi_T} \int_{|z| > r} [A(u_{\Delta y}(s, y)) - A(u_{\Delta y}(s, y + z))] \rho_\lambda(z) \psi(t, x) \rho_{\delta_0}(t - s) \left( \mathcal{P}_\lambda(x, y) - \rho_\delta(x - y) \right) \\
\times \eta'(u_{\Delta y}(s, y) - k) \, dy \, dx \, dt \right].
\]
Lemma 4.6. It holds that
\[ |C| \leq \mathbb{E} \left[ \int_{\Pi_T \times \Pi_T} \int_{|z| > r} \left| A(u_{\Delta y}(s, y)) - A(u_{\Delta y}(s, y + z)) \right| \rho_{\delta}(t-s) \left| \mathcal{P}_{\delta}^y(x, y) - \rho_{\delta}(x-y) \right| \, dy \, ds \, dx \, dt \right] \]
\[ \leq \|A'\|_{\infty} \|\psi\|_{\infty} \mathbb{E} \left[ \int_{|z| > r} \left| u_{\Delta y}(s, y) - u_{\Delta y}(s, y + z) \right| \rho_{\delta}(z) \int_{\mathbb{R}^s} \left| \mathcal{P}_{\delta}^y(x, y) - \rho_{\delta}(x-y) \right| \, dx \, dy \, ds \right] \]
\[ \leq C \frac{\Delta y}{\delta} \|A'\|_{\infty} \|\psi\|_{\infty} \left[ \|u_0\|_{BV} \int_{|z| > r} |z| \rho_{\delta}(z) \int_{|z| > 1} \rho_{\delta}(z) \, dx \, dy \, ds \right], \]
where to derive the last inequality we have used the fact that \( \int_{\mathbb{R}^s} |\mathcal{P}_{\delta}(y) - \rho_{\delta}(y)| \, dy \leq \frac{\Delta y}{\delta} \) for any space mollifier \( \rho_{\delta} \). Now we conclude the proof of the lemma with the aid of the following:
\[ \int_{|z| > \Delta y} |z| \rho_{\delta}(z) \leq \begin{cases} C_\lambda, & \text{if } \lambda < 1/2, \\ C_\lambda \ln \Delta y, & \text{if } \lambda = 1/2, \\ C_\lambda \Delta y^{1-2\lambda}, & \text{if } \lambda > 1/2. \end{cases} \]

Proof. Following [9, Lemma 3.5] we have,
\[ \lim_{\delta \to 0} \sup_{\gamma \to \infty} \lim_{\delta_0 \to 0} \int_{I_7} \leq -E \left[ \int_{\Pi_T \times \mathbb{R}^d} |A(u_t(t,x)) - A(k)| \mathcal{E}_{\lambda,r}[\psi(t, \cdot)|\rho_{\delta}(\cdot-y)](x) \zeta(u_{\Delta y}(t, y) - k) \, dk \, dy \, dx \, dt \right] \]
\[ = E \left[ \int_{\Pi_T \times \mathbb{R}^d} A(u_t(t,x)) - A(k) \int_{|z| \leq r} (\tau - 1)z^T D^2(\psi(t,x)|\rho_{\delta}(x-y+\tau z)).z \, d\tau \, d\rho_{\delta}(z) \right] \]
\[ \leq E \left[ \int_{\Pi_T \times \mathbb{R}^d} \int_{|z| \leq r} \int_0^1 |\nabla A(u_t(t,x))|\nabla(\psi(t,x)|\rho_{\delta}(x-y+\tau z)) \|\nabla \psi(t, \cdot)|_{\infty}^{2} \, d\tau \, d\rho_{\delta}(z) \right] \]
\[ + E \left[ \int_{\Pi_T \times \mathbb{R}^d} \int_{|z| \leq r} \int_0^1 |\nabla A(u_t(t,x))|\nabla(\psi(t,x)|\rho_{\delta}(x-y+\tau z)) \|\nabla \psi(t, \cdot)|_{\infty}^{2} \, d\tau \, d\rho_{\delta}(z) \right], \]
\[ \leq \|u_0\|_{BV} \|A'\|_{\infty} \int_0^T \left( \|\nabla \psi(t, \cdot)|_{\infty}^{2} + \frac{C}{\delta} \|\psi(t, \cdot)|_{\infty} \right) \int_{|z| \leq r} |z|^2 \, d\rho_{\delta}(z) \, dt. \]
To handle the other term we consider,
\[ I_7 = -E \left[ \int_{\Pi_T \times \Pi_T} A_k^y(u_{\Delta y}(s, y)) \frac{d}{\Delta y} \left[ \text{r}_\delta^y(t, x, s, \cdot) \right](y) \zeta(u_t(t, x) - k) \, dy \, ds \, dx \, dt \right] \]
\[ = -E \left[ \int_{\Pi_T \times \Pi_T} A_k^y(u_{\Delta y}(s, y)) \int_{|\mathbf{z}| \leq r} \left( \text{r}_\delta^y(x, y) - \text{r}_\delta^y(x, y + z) \right) \rho_{\delta}(z) \right. \]
\[ \times \psi(t,x)|\rho_{\delta}(t-s)| \zeta(u_t(t, x) - k) \, dy \, ds \, dx \, dt \].
Following [14], we denote in a first step
\[ I_7^0 = -E \left[ \int_{\Pi_T \times \Pi_T \times \mathbb{R}} A_0^\delta(u_{\Delta y}(s, y)) \int_{|z| \leq r} \left( \psi_0^\delta(x, y) - (\psi_0^\delta)\right) \right] d\mu_\delta(z) \]
\[ \times \psi(t, x) \rho_\delta(t - s) \delta(u_s(t, x) - k) \, dk \, dy \, ds \, dx \, dt, \]
where \((\psi_0^\delta)\) denotes a regularization of \(\psi_0^\delta\) in the \(y\)-variable by a mollification of parameter \(\theta\). Then,
\[ |I_7^0| = E \left[ \int_{\Pi_T \times \Pi_T \times \mathbb{R}} A_0^\delta(u_{\Delta y}(s, y)) \int_{|z| \leq r} \int_0^1 (1 - \tau) z^T D^2(\psi_0^\delta)\right] d\mu_\delta(z) \]
\[ \times \psi(t, x) \rho_\delta(t - s) \delta(u_s(t, x) - k) \, dk \, dy \, ds \, dx \, dt \]
\[ = E \left[ \int_{\Pi_T \times \Pi_T \times \mathbb{R}} \int_{|z| \leq r} \int_0^1 (1 - \tau) \nabla(\psi_0^\delta)\right] d\mu_\delta(z) \]
\[ \times \psi(t, x) \rho_\delta(t - s) \delta(u_s(t, x) - k) \, dk \, dy \, ds \, dx \, dt \]
\[ \leq \|A'\|_{L^\infty} E \left[ \int_{\Pi_T \times \Pi_T} \int_{|z| \leq r} \int_0^1 |\nabla(\psi_0^\delta)| d\mu_\delta(z) \right] \]
\[ \times \psi(t, x) \rho_\delta(t - s) \, dk \, dy \, ds \, dx \, dt \]
\[ \leq \frac{C \|A'\|_{L^\infty} \|u_0\|_{L^\infty} \|\psi\|_{L^\infty}}{\delta} \int_{|z| \leq r} |z|^2 \, d\mu_\delta(z) \leq \frac{C \|A'\|_{L^\infty} \|u_0\|_{L^\infty} \|\psi\|_{L^\infty}}{\delta} \int_{|z| \leq r} |z|^2 \, d\mu_\delta(z) \leq \frac{C \|A'\|_{L^\infty} \|u_0\|_{L^\infty} \|\psi\|_{L^\infty}}{\delta} , \]
where to derive the penultimate inequality we follow [14, Proof of Thm 7.1] and use the fact that 
\(\|\psi_0^\delta\|_{L^\infty} \leq \|\psi_0^\delta\|_{L^\infty} \leq \frac{C}{\delta} \). Finally to get the same estimate for \(I_7\), we use \(\lim_{\delta \to 0} I_7^0 = I_7\). □

Thus making use of Lemmas 4.1- 4.6, and recalling that \(r = \Delta y\), we pass to the limit in \(\varepsilon\) to get
\[ 0 \leq E \left[ \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \left| u_{\Delta y}(0, y) - u_0(x) \right| \, g_\delta(x - y) \psi(0, x) \, dy \, dx \right] \]
\[ + E \left[ \int_{\Pi_T} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \left| u(t, x) - u_{\Delta y}(t, y) \right| \, g_\delta(x - y) \beta_{\psi}(t, x) \, dy \, dx \, dt \right] \]
\[ + E \left[ \int_{\Pi_T} \int_{\mathbb{R}^d} Q \left( u(t, x), u_{\Delta y}(t, y) \right) \cdot \nabla \psi(t, x) \, g_\delta(x - y) \, dy \, dx \, dt \right] \]
\[ + E \left[ \int_{\Pi_T} \int_{\mathbb{R}^d} \left| \Delta y u_{\Delta y}(t, y) \right| \, \psi(t, x) \, \beta_{\psi}(t, x) \, g_\delta(x - y) \, dy \, dx \, dt \right] \]
\[ + \frac{C A}{\delta} \|u_0\|_{L^\infty} \|\psi\|_{L^\infty} \left\{ \begin{array}{ll} \Delta y, & \text{if } \lambda < 1/2 \\ \Delta y \ln |\Delta y|, & \text{if } \lambda = 1/2 \\ \Delta y^{2(1-\lambda)}, & \text{if } \lambda > 1/2 \end{array} \right\} \int_{|z| \leq \Delta y} |z|^2 \, d\mu_\delta(z) \, dt + \frac{C}{\delta} \left( \frac{\Delta y}{\delta} \right)^{2-2\lambda} \|\psi\|_{L^\infty} \int_{|z| \leq \Delta y} |z|^2 \, d\mu_\delta(z) \, dt \].

To proceed further, we make a special choice for the function \(\psi(t, x)\). To this end, for each \(h > 0\) and fixed \(t \geq 0\), we define
\[ \psi_h^t(s) = \begin{cases} 1, & \text{if } s \leq t, \\
1 - \frac{s - t}{h}, & \text{if } t \leq s \leq t + h, \\
0, & \text{if } s \geq t + h. \end{cases} \]
Furthermore, let $\rho$ be any non-negative mollifier. Clearly, (4.6) holds with $\psi(s,x) = \psi_R^\rho(s) (\psi_R * \rho)(x)$. At this point, we can closely follow Bhauryal et al. [8, 9] and pass to the limits as $R \to \infty$ and $h \to 0$ to conclude

$$
\mathbb{E} \left[ \int_{\mathbb{R}^2} |u(t, y) - u_{\Delta y}(t, y)| \, dx \right] \leq C \mathbb{E} \left[ \int_{\mathbb{R}^2} |u_{\Delta y}(0, y) - u_0(y)| \, dx \right] + C \left( \delta + \frac{(\Delta y)^2 - 2\lambda}{\delta} \right) + C \left\{ \begin{array}{ll}
\delta, & \text{if } \lambda < 1/2, \\
\Delta y \ln \Delta y, & \text{if } \lambda = 1/2, \\
\Delta y^{2(1-\lambda)}, & \text{if } \lambda > 1/2.
\end{array} \right.
$$

For $\lambda \in (0, 1/2]$, we choose $\delta = \sqrt{\Delta y}$ and for $\lambda \in (1/2, 1)$, we choose $\delta = (\Delta y)^{1-\lambda}$ to obtain the following

$$
\mathbb{E} \left[ \int_{\mathbb{R}^2} |u(t, y) - u_{\Delta y}(t, y)| \, dy \right] \leq C \left\{ \begin{array}{ll}
\sqrt{\Delta y}, & \text{if } \lambda < 1/2, \\
\sqrt{\Delta y} \ln \Delta y, & \text{if } \lambda = 1/2, \\
\Delta y^{(1-\lambda)}, & \text{if } \lambda > 1/2,
\end{array} \right.
$$

provided the initial error satisfies $\mathbb{E} \left[ \int_{\mathbb{R}^2} |u_{\Delta y}(0, y) - u_0(y)| \, dy \right] \leq \sqrt{\Delta y}$.

5. Numerical Experiments

In this section, we simulate numerical experiments to substantiate the results we have shown in the previous sections. In what follows, inspired by Del Teso et al. [19, 20] for the deterministic fractional porous medium operator, we test numerically the performance of the proposed scheme (2.5). Here we use a Godunov scheme for the first-order operator and an explicit scheme for the noise term. We set the underlying target equation (1.1) (posed in $\mathbb{R}$) with

$$
A(u) = (u - \frac{1}{2})^+, \quad f(u) = \frac{1}{2} u^2, \quad \sigma(u) = u(1 - u), \quad u_0(x) = 2e^{-x^2/2} \mathbb{1}_{(-1, 1)}.
$$

Let us mention that this configuration of data leads to a solution $u$ satisfying $0 \leq u \leq 1$ so that one may replace $f(u)$ and $\sigma(u)$ by $\frac{1}{2} \min(1, |u|)^2$ and $u^+(1 - u)^+$ respectively to be compatible with the assumptions of the paper.

Next, we present the methodology of the fully discrete explicit numerical scheme. $\Delta t > 0$ and $\Delta x > 0$ are the time-step and the spatial mesh size respectively. $t^n = n\Delta t$ for $n = 0, 1, 2, \cdots, N = \frac{T}{\Delta t}$ denotes the temporal grid and $x_i = i\Delta x$ for $i \in \mathbb{Z}$ the spatial one. The scheme becomes

$$
U_{i}^{n+1} = U_i^n \Delta x \frac{d}{\Delta x} E(U_i^n, U_{i+1}^n) - \frac{\Delta t}{\Delta x} \sum_{j \in \mathbb{Z}} G_{i-j} A(U_i^n) + \sigma(U_i^n) \left( W((t_{n+1})_{\Delta t}) - W(t_{n\Delta t}) \right),
$$

(5.1)

$$
U_i^0 = \frac{1}{\Delta x} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} u_0(x) \, dx,
$$

(5.2)

where $U_i^n$ is the approximate solution of (1.1) in the cell $[t^n, t^{n+1}] \times [x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}]$ and we denote the weights by $G_{i-j} := G_{i+j}$ for $i \neq j$, and $G_0 := G_{i,j}$. The numerical solution is the piecewise constant function denoted by

$$
u_{\Delta x}(t, x) = U_i^n, \quad \text{for all } (t, x) \in [t^n, t^{n+1}] \times [x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}].
$$

A tedious but straightforward calculation reveals the explicit values of the weights:

| $i$ | $\frac{G_i}{\Delta x}$ | $\lambda \neq \frac{1}{2}$ | $\lambda = \frac{1}{2}$ |
|-----|---------------------|------------------|------------------|
| $|i| \geq 2$ | $- \frac{d_1((\Delta x)^{2-\lambda})}{2\lambda(1-2\lambda)} \left[ (i-1)^{1-2\lambda} + 2(i)^{1-2\lambda} + (i+1)^{1-2\lambda} \right]$ | $\frac{d_1}{\Delta x} \ln \left[ \frac{i^2}{i^2-1} \right]$ |
| $|i| = 1$ | $\frac{d_1 \Delta x^{2-\lambda}}{2\lambda(1-2\lambda)} \left[ 2^{-1-2\lambda} + 2\lambda - 2 \right] = \frac{d_1 \Delta x^{2-\lambda}}{2\lambda(1-2\lambda)} \left[ 2^{-1-2\lambda} + \lambda \Delta x^{2\lambda} - 2 \right]$ | $- \frac{d_1 \Delta x^{2-\lambda}}{2\lambda(1-2\lambda)}$ |
| $i = 0$ | $\frac{d_1((\Delta x)^{2-\lambda})}{\lambda(1-2\lambda)} \left[ (1)^{1-2\lambda} + 2(i)^{1-2\lambda} \right] = \frac{d_1((\Delta x)^{2-\lambda})}{\lambda(1-2\lambda)} \left[ 1 - \lambda \Delta x^{2\lambda} \right]$ | $\frac{d_1 \Delta x^{2-\lambda}}{2\lambda(1-2\lambda)}$ |

Here $d_1 = \frac{2\Gamma(\frac{1+\lambda}{2})}{\pi^{1/2} \Gamma(1-\lambda)}$, with $\Gamma$ being the classical Gamma function.
5.1. **Computation of the non-local diffusion term:** Following [21], a truncated domain \([-K \Delta x, K \Delta x]\) is considered for a given \(K\), and one considers that \(U_i = U_{-K}\) for all \(i \leq -K\) and \(U_i = U_K\) for all \(i \geq K\). Therefore, the non-local term in (5.1) becomes,

\[
\frac{\Delta t}{\Delta x} \sum_{j \in \mathbb{Z}} \tilde{G}_{j-i} A(U^n_j) = \frac{\Delta t}{\Delta x} \sum_{|j| < K} \tilde{G}_{j-i} A(U^n_j) + \frac{\Delta t}{\Delta x} A(U^n_{-K}) \sum_{j \leq -K} \tilde{G}_{j} + \frac{\Delta t}{\Delta x} A(U^n_K) \sum_{j \geq K} \tilde{G}_{j} \\
= \frac{\Delta t}{\Delta x} \left[ \sum_{|j| < K} \tilde{G}_{j-i} A(U^n_j) - \frac{1}{2} A(U^n_{-K}) \sum_{|j| < K+i} \tilde{G}_{j} - \frac{1}{2} A(U^n_K) \sum_{|j| < K-i} \tilde{G}_{j} \right]
\]

since for any \(B > 0\), \(\sum_{j \leq -B} \tilde{G}_{j} = \sum_{j \geq B} \tilde{G}_{j} = \frac{1}{2} \sum_{|j| \geq B} \tilde{G}_{j} = -\frac{1}{2} \sum_{|j| < B} \tilde{G}_{j}\).

5.2. **Numerical Examples:** We have chosen to work with the truncated spatial domain \([-3, 3]\) and the time of simulation is \(T = 1\). We considered as a very small time step \(\Delta t = 2^{-12}\), five times steps \(\Delta t \in \{2^{-9}, 2^{-8}, 2^{-7}, 2^{-6}, 2^{-5}\}\) and the corresponding space steps: \(\Delta x\) and five \(\Delta x\), given by using a CFL condition. The later is based on the classical one for monotone flux: \(\|f'\|_{\infty} \Delta x \sim 1\) for \(\lambda \leq 0.5\) and on the power \(\lambda\) and the weights: \(\tilde{G}_0 \Delta x \sim 1\) else.

Since one doesn’t know about explicit solutions for such problems, one proposes as a numerical rate of convergence:

\[
\text{Error} := \max_{t \in \{0.25, 0.5, 0.75, 1\}} \mathbb{E} \left[ \|u_{\Delta x}(t) - u_{\Delta x}(t)\|_{L^1(-3,3)} \right],
\]

where \(\mathbb{E}\) denotes the statistical average over 5000 independent paths.

This “Error” is then calculated on the 1088 processing cores research computing cluster “Pyrene” (univ. Pau) for five values of \(\lambda\) in \(\{0.8, 0.65, 0.5, 0.3, 0.1\}\) and the corresponding results are given in the figures below. One can note that the numerical rate of convergence seems to be of order 1 for smaller values of \(\lambda\), and it is of order 1/2 for larger values of \(\lambda\).

| \(\Delta x\)   | Error | Rate |
|---------------|-------|------|
| 0.78 E-2      | 0.52 E-2 |      |
| 1.56 E-2      | 1.07 E-2 | 1.02 |
| 3.12 E-2      | 2.05 E-2 | 0.93 |
| 6.25 E-2      | 3.81 E-2 | 0.89 |
| 12.5 E-2      | 6.97 E-2 | 0.88 |

Table 1. \(\lambda = 0.1\)

| \(\Delta x\)   | Error | Rate |
|---------------|-------|------|
| 0.78 E-2      | 0.53 E-2 |      |
| 1.56 E-2      | 1.09 E-2 | 1.03 |
| 3.12 E-2      | 2.12 E-2 | 0.95 |
| 6.25 E-2      | 4.01 E-2 | 0.92 |
| 12.5 E-2      | 7.44 E-2 | 0.88 |

Table 2. \(\lambda = 0.3\)

| \(\Delta x\)   | Error | Rate |
|---------------|-------|------|
| 0.78 E-2      | 0.50 E-2 |      |
| 1.56 E-2      | 1.04 E-2 | 1.05 |
| 3.12 E-2      | 2.04 E-2 | 0.96 |
| 6.25 E-2      | 3.90 E-2 | 0.99 |
| 12.5 E-2      | 7.34 E-2 | 0.90 |

Table 3. \(\lambda = 0.5\)
## Table 4. $\lambda = 0.65$

| $\Delta x$ | Error Rate |
|------------|-------------|
| 0.78 E-2   | 0.68 E-2    |
| 1.56 E-2   | 1.30 E-2    |
| 3.12 E-2   | 2.24 E-2    |
| 6.25 E-2   | 4.01 E-2    |
| 12.5 E-2   | 6.89 E-2    |

## Table 5. $\lambda = 0.8$

| $\Delta x$ | Error Rate |
|------------|-------------|
| 0.78 E-2   | 0.92 E-2    |
| 1.56 E-2   | 1.52 E-2    |
| 3.12 E-2   | 2.45 E-2    |
| 6.25 E-2   | 3.62 E-2    |
| 12.5 E-2   | 5.12 E-2    |

Tables 1 to 5: Information about the numerical rate of convergence.

## References

1. N. Alibaud. Entropy formulation for fractal conservation laws. *J. Evol. Equ.*, 7(1), 145-175, 2007. 1.1
2. C. Bauzet, G. Vallet and P. Wittbold. The Cauchy problem for a conservation law with a multiplicative stochastic perturbation. *Journal of Hyperbolic Differential Equations*, 2012. 1.1, 4.1, 4.4
3. C. Bauzet, G. Vallet and P. Wittbold. A degenerate parabolic-hyperbolic Cauchy problem with a stochastic force. *Journal of Hyperbolic Differential Equations*, 12(3) (2015) 501-533. 1.1
4. C. Bauzet, G. Vallet and P. Wittbold. The Dirichlet problem for a conservation law with a multiplicative stochastic perturbation. *J. Funct. Anal.*, 266 (4), 2503-2545, 2014. 1.1
5. C. Bauzet. Time-splitting approximation of the Cauchy problem for a stochastic conservation law. *Math. Comput. Simulation* 118 (2015), 73-86. 1.1
6. C. Bauzet, J. Charrier, and T. Gallouët. Convergence of flux-splitting finite volume schemes for hyperbolic scalar conservation laws with a multiplicative stochastic perturbation. *Mathematics of Computation* 85 (2016), 2777-2813. 1.1
7. C. Bauzet, J. Charrier, and T. Gallouët. Convergence of monotone finite volume schemes for hyperbolic scalar conservation laws with multiplicative noise. *Stochastic partial differential equations: analysis and computations* 4(1) (2016), 150-223. 1.1
8. N. Bhauryal, U. Koley, G. Vallet. The Cauchy problem for a fractional conservation law driven by Lévy noise. *Stochastic Processes and their applications* https://doi.org/10.1016/j.spa.2020.03.009. 1.1
9. N. Bhauryal, U. Koley, G. Vallet. A Fractional degenerate parabolic-hyperbolic Cauchy problem with noise. https://arxiv.org/abs/2008.03141. 1.1, 2.1, 2.3, 4.2, 4.4
10. I. H. Biswas, U. Koley, and A. K. Majee. Continuous dependence estimate for conservation laws with Lévy noise. *J. Differ. Equ.*, 259 (2015), 4683-4706. 1.1
11. A. Chaudhary, and U. Koley: A convergent finite volume scheme for stochastic compressible barotropic Euler equations, Submitted, https://arxiv.org/submit/3901170. 1.1
12. A. Chaudhary, and U. Koley: On weak-strong uniqueness for stochastic equations of incompressible fluid flow, https://arxiv.org/pdf/2012.10175.pdf. 1.1
13. G. Q. Chen, Q. Ding, and K. H. Karlsen. On nonlinear stochastic balance laws. *Arch. Rational Mech. Anal.*, 204 (3), 707-743, 2012. 1.1
14. Cifani, S., Jakobsen. E.R. On numerical methods and error estimates for degenerate fractional convection diffusion equations. *Numer. Math.*, 127, 447-483 (2014) 1.1, 2.1, 2.2, 2.2.1, 4, 5.2
15. S. Cifani, and E. R. Jakobsen. Entropy solution theory for fractional degenerate convection-diffusion equations. *Ann. I. H. Poincaré*, 28(3), 413-441, 2011. 1.1
16. M. G. Crandall and A. Majda. Monotone difference approximations for scalar conservation laws. *Math. Comp.*, 34(149):1–21, 1980. 1.1
17. A. Debussche, J. Vovelle: *Scalar conservation laws with stochastic forcing*. *J. Funct. Anal.* 259(4), 1014–1042 (2010). 1.1
18. A. Debussche, M. Hofmanová, J. Vovelle, Degenerate parabolic stochastic partial differential equations: quasilinear case. *Ann. Probab.* 44 (2016), no. 3, 1916–1955. 60H15 (35K65 35R60). 1.1
19. F. Del Teso, J. Endal, and E. R. Jakobsen. Robust numerical methods for nonlocal (and local) equations of porous medium type. part II: Schemes and experiments. *SIAM Journal on Numerical Analysis*, 56(6):3611-3647, 2018. 5
20. F. Del Teso, J. Endal, and E. R. Jakobsen. Robust numerical methods for nonlocal (and local) equations of porous medium type. part I: Theory. *SIAM Journal on Numerical Analysis*, 57(5):2266-2299, 2019. 5
21. Jérôme Droniou. A numerical method for fractional conservation laws. *Mathematics of Computation*, 269, 95–124, 2010. 1.1, 5.1, 5.2
22. R. Eymard, T. Gallouët, and R. Herbin. Finite volume methods. *Handbook of numerical analysis*, Vol. VII, *Handb. Numer. Anal.*, VII,713-1020. North-Holland, Amsterdam, 2000.
23. A. Harten, P. D. Lax, and B. van Leer. On upstream differencing and Godunov-type schemes for hyperbolic conservation laws. *SIAM Rev.*, 25(1):35–61, 1983. 1.1
24. M. Hofmanova, U. Koley, and U. Sarkar: Measure-valued solutions to the stochastic compressible Euler equations and incompressible limits. https://arxiv.org/pdf/2012.07391.pdf. 1.1
25. H. Holden and N. H. Risebro. Conservation laws with random source. *Appl. Math. Optim.*, 36(1997), 229-241.
[26] Y. Huang and A. Oberman. Numerical methods for the fractional laplacian: A finite difference-quadrature approach. *SIAM Journal on Numerical Analysis*, 52(6):3056–3084, 2014. 5.2

[27] K. H. Karlsen, U. Koley, and N. H. Risebro An error estimate for the finite difference approximation to degenerate convection-diffusion equations. *Numer. Math.*, 121(2): 367-395, 2012. 1.1

[28] U. Koley, A. K. Majee, and G. Vallet. Continuous dependence estimate for a degenerate parabolic-hyperbolic equation with Lévy noise. *Stochastic Partial Differential Equations: Analysis and Computations*, to appear, DOI: 10.1007/s40072-016-0084-z 1.1

[29] U. Koley, A. K. Majee, and G. Vallet. A finite difference scheme for conservation laws driven by Lévy noise. *IMA Journal of Numerical Analysis*, 38(2), 998-1050, 2018 1.1

[30] U. Koley, N. H. Risebro, C. Schwab and F. Weber. A multilevel Monte Carlo finite difference method for random scalar degenerate convection-diffusion equations. *J. Hyperbolic Differ. Equ.*, 14(3), 415-454, 2017. 1.1

[31] U. Koley, D. Ray, and T. Sarkar. Multi-level Monte Carlo finite difference methods for fractional conservation laws with random data., *SIAM/ASA J. Uncert., Quantif.*, 9(1), 65–105, 2021. 1.1

[32] I. Kroker and C. Rohde. Finite volume schemes for hyperbolic balance laws with multiplicative noise. *Applied Numerical Mathematics* 62, 441-456, 2012. 1.1

[33] S. N. Kruzkov. First order quasilinear equations with several independent variables. *Mat. Sb. (N.S.)*, 81(123): 228-255, 1970. 1.1

[34] O. A. Oleinik. Convergence of certain difference schemes. *Soviet Math. Dokl.*, 2:313–316, 1961. 1.1

[35] Prévôt, Claudia and Röckner, Michael *A concise course on stochastic partial differential equations*. volume 1905 of *Lecture Notes in Mathematics*. Springer, 2007 2, 2.2.1

[36] A. I. Vol’pert. Generalized solutions of degenerate second-order quasilinear parabolic and elliptic equations. *Adv. Differential Equations*, 5(10-12):1493–1518, 2000. 1.1

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