Observation of the $\Xi_0^{0}$ Baryon

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**Abstract**

The observation of the bottom, strange baryon $\Xi^0_b$ through the decay chain $\Xi^0_b \rightarrow \Xi^+_c \pi^-$, where $\Xi^+_c \rightarrow \Xi^- \pi^+ \pi^+$, $\Xi^- \rightarrow \Lambda \pi^-$, and $\Lambda \rightarrow p \pi^-$, is reported using data corresponding to an integrated luminosity of 4.2 fb$^{-1}$ from $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV recorded with the Collider Detector at Fermilab. A signal of $25.3^{+5.6}_{-5.4} \pm 6.6$ candidates is observed whose probability of arising from a background fluctuation is $3.6 \times 10^{-12}$, corresponding to 6.8 Gaussian standard deviations. The $\Xi^0_b$ mass is measured to be $5787.8 \pm 5.0$ (stat) $\pm 1.3$ (syst) MeV/$c^2$. In addition, the $\Xi^-_b$ baryon is observed through the process $\Xi^-_b \rightarrow \Xi^0_c \pi^-$, where $\Xi^0_c \rightarrow \Xi^- \pi^+$, $\Xi^- \rightarrow \Lambda \pi^-$, and $\Lambda \rightarrow p \pi^-$. 

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The quark model has had great success in describing the spectroscopy of hadrons. For the $c$ and $b$ mesons, all of the ground states have been observed \[1\]. The spectroscopy of $c$ baryons also agrees well with the quark model, and a rich spectrum of baryons containing $b$ quarks is predicted \[2\]. Until recently, direct observation of $b$ baryons has been limited to a single state, the $\Lambda^0_b$ (quark content $|udb\rangle$) \[1\]. The accumulation of large data sets from the Tevatron has improved this situation and made possible the observation of the $\Xi_b^-$ ($|dsb\rangle$) \[3, 4\], the $\Sigma^{(*)}_b$ states ($|uub\rangle, |ddb\rangle$) \[5\], and the $\Omega_b$ ($|ssb\rangle$) \[6, 7\].

In this paper, we report the observation of an additional heavy baryon and the measurement of its mass. The decay properties of this state are consistent with the weak decay of a $b$ baryon. We interpret the result as the observation of the $\Xi^0_b$ baryon ($|usb\rangle$). This measurement is made in $p\bar{p}$ collisions at a center of mass energy of 1.96 TeV using the Collider Detector at Fermilab (CDF II), by fully reconstructing the decay chain $\Xi^0_b \to \Xi^+_c \pi^-$, where $\Xi^+_c \to \Xi^- \pi^+ \pi^+$, $\Xi^- \to \Lambda \pi^-$, and $\Lambda \to p \pi^-$. Charge conjugate modes are included implicitly. In addition, we observe the $\Xi_b^-$ through the similar decay chain $\Xi_b^- \to \Xi^0_c \pi^-$, where $\Xi^0_c \to \Xi^- \pi^+$, $\Xi^- \to \Lambda \pi^-$, and $\Lambda \to p \pi^-$. These studies use a data sample corresponding to an integrated luminosity of 4.2 fb$^{-1}$ and constitute the first exclusive reconstruction of the $\Xi^0_b$ and the first for the $\Xi_b^-$ in this decay channel.

The CDF II detector has been described in detail elsewhere \[8\]. This analysis relies upon the tracking system that operates inside a 1.4 T solenoidal magnetic field. A five-layer silicon detector (SVX II) measures track positions at radii of 2.5 to 10.6 cm to provide high precision impact parameter measurements. Each of these layers provides a transverse measurement and a stereo measurement of 90° (three layers) or $\pm 1.2^\circ$ (two layers) with respect to the beam direction. An open-cell drift chamber (COT) covers the radial region from 43 cm to 132 cm and provides track momentum measurement.

Data acquisition is triggered by a system designed to collect particle candidates that decay with lifetimes characteristic of heavy flavor hadrons. The first level of the trigger system requires two tracks in the COT with transverse momentum $p_T > 2.0$ GeV/c. In the second level of the trigger, the silicon vertex trigger \[9\] is used to associate SVX II data with the tracks found in the COT and provides precise impact parameter resolution (typically 40 $\mu$m) for these tracks. The silicon vertex trigger requires two tracks with impact parameters in the range 0.1-1.0 mm with respect to the beam and a point of intersection that is measured with at least a 200 $\mu$m displacement transverse to the beam.
This analysis combines the trajectories of charged particles to infer the presence of several different hadrons in the decay chains. The decay point for each weak decay process in the decay chain is reconstructed and used to identify the corresponding hadron. Consequently, it is useful to define two quantities that are used frequently throughout the analysis that relate the paths of weakly decaying objects to their points of origin. Both quantities are defined in the transverse view and make use of the point of closest approach $\vec{r}_c$ of the particle trajectory to a point of origin $\vec{r}_o$ and of the measured particle decay position $\vec{r}_d$. The first quantity used here is transverse flight distance $f(h)$ of hadron $h$, which is the distance a particle has traveled in the transverse view. For neutral particles, flight distance is given by $f(h) \equiv (\vec{r}_d - \vec{r}_o) \cdot \vec{p}_T(h)/|\vec{p}_T(h)|$, where $\vec{p}_T(h)$ is the transverse momentum of the hadron candidate. For charged particles, the flight distance is calculated as the arc length in the transverse view from $\vec{r}_c$ to $\vec{r}_d$. Flight distance is used to calculate the proper decay time of weakly decaying states, where the decay time is given by $t \equiv f(h)M(h)/(c|\vec{p}_T(h)|)$, where $M(h)$ is the reconstructed mass.

A complementary quantity used in this analysis is transverse impact distance $d(h)$, which is the distance of the point of closest approach to the point of origin. For neutral particles, transverse impact distance is given by $d(h) \equiv |(\vec{r}_d - \vec{r}_o) \times \vec{p}_T(h)|/|\vec{p}_T(h)|$. The impact distance of charged particles is simply the distance from $\vec{r}_c$ to the point of origin.

The reconstruction of $\Lambda$ candidates uses all tracks with $p_T > 0.4$ GeV/c found in the COT. Pairs of oppositely charged tracks are combined to identify these neutral decay candidates, and silicon detector information is not used due to the large transverse displacement of the $\Lambda$ decay. Candidate selection is based upon the mass calculated for each track pair, which has a resolution of 1.5-2.0 MeV/c$^2$ and is required to fall within 9 MeV/c$^2$ of the nominal $\Lambda$ mass [1] after the appropriate mass assignment for each track. The proton (pion) mass is assigned to the track with the higher (lower) momentum. This mass assignment is always correct for the $\Lambda$ candidates used in this analysis because of the kinematics of $\Lambda$ decay and the lower limit in the transverse momentum acceptance of the tracking system. Background to the $\Lambda$ ($c\tau = 7.9$ cm) [1] is reduced by requiring the transverse flight distance of the $\Lambda$ from the beam position to be greater than 1.0 cm, which corresponds to typically 0.6 $\sigma_f$, where $\sigma_f$ is the flight distance resolution.

For events that contain a $\Lambda$ candidate, the remaining tracks reconstructed in the COT, again without additional silicon information, are assigned the pion mass, and $\Lambda \pi^-$ combina-
A search for weak decays of the hyperons \( \Xi^- \) is conducted, focusing on the decay process \( \Xi^- \rightarrow \Lambda \pi^- \). Several features of the track topology are used to reduce the background to this process. In order to obtain the best possible mass resolution for \( \Xi^- \) candidates, the reconstruction requires a convergent fit of the three tracks that simultaneously constrains the \( \Lambda \) decay products to the \( \Lambda \) mass and the \( \Lambda \) trajectory to intersect with the helix of the \( \pi^- \) originating from the \( \Xi^- \) candidate. The \( \Lambda \pi^- \) mass obtained from this fit has a resolution comparable to the \( \Lambda \) and is required to fall within 9 MeV/c\(^2\) of the nominal \( \Xi^- \) mass [1]. In addition, the flight distance of the \( \Lambda \) candidate with respect to the reconstructed decay point of the \( \Xi^- \) candidate is required to exceed 1.0 cm. Similarly, due to the long lifetime of the weakly decaying \( \Xi^- \) (\( c\tau = 4.9 \) cm) [1], a transverse flight distance of at least 1.0 cm (which typically corresponds to 1.0 \( \sigma_f \)) with respect to the beam position is required.

In some instances, the intersection of the \( \pi^- \) helix with the \( \Lambda \) trajectory produces a situation where two \( \Lambda \pi^- \) vertices satisfy the constrained fit and displacement requirements. In addition, the complexity of the \( \Xi^- \) and \( \Lambda \) decays allows for occasional combinations where the proper identity of the three tracks is ambiguous. A single, preferred candidate is chosen by retaining only the fit combination with the highest probability of satisfying the constrained fit.

The kinematics of hyperon decay and the lower \( p_T \) limit of 0.4 GeV/c on the decay daughter tracks force the majority of \( \Xi^- \) candidates to have \( p_T > 1.5 \) GeV/c. This fact, along with the long lifetime of the \( \Xi^- \), results in a significant fraction of the hyperon candidates having decay points located several centimeters radially outward from the beam position. Therefore, we are able to refine the \( \Xi^- \) reconstruction by making use of the improved determination of the trajectory that can be obtained by tracking these particles in the silicon detector. The \( \Xi^- \) candidates have an additional fit performed with the three tracks that simultaneously constrains both the \( \Lambda \) and \( \Xi^- \) masses of the appropriate track combinations and provides the best possible estimate of the hyperon momentum and decay position. The result of this fit is used to define a helix that serves as the seed for an algorithm that associates silicon detector hits with the \( \Xi^- \) track. Candidates with track measurements in at least one layer of the silicon detector have excellent impact distance resolution (typically 60 \( \mu \)m).

The samples of \( \Xi_c^- \) and \( \Xi_c^+ \) candidates used in this analysis are obtained by combining the \( \Xi^- \) candidates that have SVX II information with additional \( \pi^+ \) candidates. The \( \pi^+ \)}
candidates are tracks that have been reconstructed with data from at least three SVX II layers. The \( \pi^+ \) used for the \( \Xi^0_c \) reconstruction is required to be consistent with the trigger requirements. The \( \Xi^+_c \) candidates are required to have at least one \( \pi^+ \) track consistent with the trigger requirements. All \( \Xi^- \pi^+(\pi^+) \) combinations are required to satisfy a constrained fit that includes mass constraints on the \( \Lambda \) and \( \Xi^- \) candidates. The mass distributions of the combinations that also satisfy \( ct > 100 \mu m \) and \( p_T > 4.0 \text{ GeV/c} \) requirements are shown in Fig. 1. Candidates with a reconstructed mass within 30(25) MeV/c\(^2\) of the nominal \( \Xi^0_c(\Xi^+_c) \) mass are used for \( b \) baryon reconstruction.

The \( \Xi_b^{(-,0)} \) candidates are reconstructed by combining the \( \Xi_c^{(0,+)} \) candidates with \( \pi^- \) candidates that satisfy the trigger requirements. The \( \Xi_b \) candidates are required to have \( p_T > 6.0 \text{ GeV/c} \), restricting the sample to candidates that are within the kinematic range where our acceptance is well modeled \[7\]. All \( \Xi_c \pi^- \) combinations are required to satisfy a constrained fit that includes mass constraints on the \( \Lambda, \Xi^-, \) and \( \Xi_c \) candidates. Combinations that are inconsistent with having originated from the collision are rejected by imposing an upper limit on the impact distance \( d_{PV} \) of the \( \Xi_b \) candidate measured with respect to the primary vertex. In addition, the full reconstruction of the \( \Xi_b \) decay chain provides an opportunity to impose a requirement on the decay time of the \( \Xi_c \) candidate since both its point of creation and decay are reconstructed.

The mean life of the charm baryons varies over a wide range and is large compared to the typical decay time resolution of 20 - 60 \( \mu m/c \) that we measure. Therefore, we have chosen a selection on the \( \Xi_c \) decay time that uses the decay time resolution \( \sigma_t \) calculated for each candidate and the mean life of the decaying state. The selection is developed by using \( \Lambda^0_b \) as a reference signal. A sample of \( \Lambda^0_b \rightarrow \Lambda^+_c \pi^- \) candidates \[10\] is used to optimize selection criteria for \( \Lambda^+_c \) decay time based on the mean life of the \( \Lambda^+_c \) and its decay time resolution. As a result of this study, we require that the measured decay time of the \( \Xi_c \) candidate falls within the range \(-2\sigma_t < t < 3\tau + 2\sigma_t \) where \( \tau \) is the mean life of the \( \Xi^0_c(cT = 33 \mu m) \) and \( \Xi^+_c(cT = 132 \mu m) \) candidates. This requirement is found to be approximately 95% efficient on our \( \Lambda^0_b \) (\( cT = 60 \mu m \)) sample and to reduce the background substantially.

The \( \Xi^0_c \pi^- \) and \( \Xi^+_c \pi^- \) mass distributions with \( d_{PV} < 100 \mu m \) and \( ct > 100 \mu m \) are shown in Fig. 2. These distributions show clear evidence of an excess near a mass of 5.8 GeV/c\(^2\) with a width consistent with our expected mass measurement resolution. The mass, yield, and significance of the \( \Xi_b^{(-,0)} \) signals are obtained by performing an unbinned likelihood
fit on the mass distribution of candidates. The likelihood function that is maximized has
the form \( \mathcal{L} = \prod_i^N (f_s G(m_i, m_0, s_m \sigma_i^m) + (1 - f_s)(a_0 + a_1 m_i)) \), where \( N \) is the number of
candidates in the sample, \( G(m_i, m_0, s_m \sigma_i^m) \) is a Gaussian distribution with average \( m_0 \) and
characteristic width \( s_m \sigma_i^m \) to describe the signal, \( m_i \) is the mass obtained for a single \( \Xi(0,+) \pi^- \) candidate, \( \sigma_i^m \) is the calculated uncertainty on
\( m_i \), and the \( a_n \) terms model the background. The quantities obtained from the fitting procedure include the fraction \( f_s \) of the candidates identified as signal, the best average mass value \( m_0 \), a scale factor on the mass resolution
\( s_m \) to allow for inaccuracy of the resolution estimate, and the values of \( a_0 \) and \( a_1 \).

For this data sample, several variations of the fit were used to test the significance. The
first of these fits corresponds to the null signal hypothesis, and fixes \( f_s = 0.0 \), \( s_m = 1.0 \),
and \( m_0 \) to the nominal mass of the \( \Xi^- \). Additional applications allow \( f_s \) to float, retain the
constraints on \( s_m \), and fix \( m_0 \) to values within 5 MeV/\( c^2 \) of the nominal mass of the \( \Xi^- \). The
value of \( -2 \ln \mathcal{L} \) for the null hypothesis exceeds the values for the fits with variable \( f_s \) by at
least 48.2 units for the \( \Xi^- \) candidate sample and by 48.3 units for the \( \Xi^0 \) candidate sample. We
interpret these as equivalent to a \( \chi^2 \) with one degree of freedom whose probability of
occurrence is \( 3.9 \times 10^{-12} \) and \( 3.6 \times 10^{-12} \), corresponding to a significance that exceeds 6.8\( \sigma \)
for both the \( \Xi^- \) and \( \Xi^0 \). We therefore interpret these results as observations of the processes
\( \Xi^- \rightarrow \Xi_c^0 \pi^- \) and \( \Xi^0 \rightarrow \Xi_c^+ \pi^- \).

Masses are obtained from the unbinned likelihood fit with the mass and resolution parameters allowed to vary. In addition, the mass fit was used on the \( \Xi^- \pi^+ \) and \( \Xi^- \pi^+ \pi^+ \) to obtain mass measurements for the \( \Xi_c^0 \) and \( \Xi_c^+ \), which are seen to be consistent with the
nominal values \([1]\). The results of these fits are listed in Table I.

| Resonance | Yield | Mass (MeV/\( c^2 \)) | Resolution Scale |
|-----------|-------|---------------------|-----------------|
| \( \Xi^0 \) | 2110 ± 70 | 2470.4 ± 0.3 | 1.16 ± 0.04 |
| \( \Xi^+ \) | 3048 ± 67 | 2467.3 ± 0.2 | 1.24 ± 0.03 |
| \( \Xi^- \) | 25.8±5.5 | 5796.7 ± 5.1 | 1.3 ± 0.2 |
| \( \Xi^0_b \) | 25.3±5.6 | 5787.8 ± 5.0 | 1.2 ± 0.2 |

The accuracy of our mass measurement scale is established by our measurements of the
\( J/\psi \), \( \psi(2S) \), and \( \Upsilon \) masses. These calibration points imply an accuracy of 0.5 MeV/\( c^2 \) on
the mass measurements of the \( \Xi^- \) and \( \Xi^- \). Our fitting technique finds that our estimate of
the mass resolution on each candidate is low, as listed in Table I. Fits where this scale factor
was fixed at 1.0 or 1.4 introduced shifts in our $\Xi^0_b$ mass result by as much as 1.0 MeV/c². A
fit with a fixed 20 MeV/c² Gaussian width, as implied by the simulation, introduced a shift
of only 0.2 MeV/c². These effects are added in quadrature with the larger of the asymmetric
nominal $\Xi_c^{(0, +)}$ mass uncertainties [1] to yield systematic uncertainties of 1.4 GeV/c² for the
$\Xi_b^-$ and 1.3 GeV/c² for the $\Xi^0_b$ mass measurements.

The momentum scale uncertainty is common to all of our mass measurements, and can
be dropped as a systematic uncertainty of a measurement of the mass difference between
the $\Xi_b^-$ and $\Xi^0_b$. Our best $\Xi_b^-$ mass measurement of 5790.9 ± 2.6(stat) ± 0.8(syst) MeV/c²
[7] is obtained from the $J/\psi\Xi^-$ final state and has a systematic uncertainty that would
be reduced to 0.6 MeV/c² without this effect. Therefore, we measure the mass difference
$M(\Xi^0_{-b}) - M(\Xi^0_{b}) = 3.1 ± 5.6$(stat) ± 1.3(syst) MeV/c², where the statistical and systematic
uncertainties of the individual measurements have been added in quadrature.

In conclusion, we have analyzed data collected with the CDF II detector at the Tevatron
to observe the bottom, strange baryon $\Xi^0_b$. The reconstruction technique is used on the $\Xi^0_b$
as well, and the observation of this state provides a cross check for the analysis. A signal
of 25.3^{+5.6}_{-5.4} $\Xi^0_b$ candidates, with a significance greater than 6σ, is seen in the decay channel
$\Xi^0_{b} \rightarrow \Xi^+_c \pi^-$ where $\Xi^+_c \rightarrow \Xi^- \pi^+ \pi^+, \Xi^- \rightarrow \Lambda \pi^-, \Lambda \rightarrow p \pi^-$. The mass of this baryon is
measured to be 5787.8 ± 5.0(stat) ± 1.3(syst) MeV/c², which is consistent with theoretical
expectations [2]. In addition, we observe 25.8^{+5.5}_{-5.2} candidates in the process $\Xi^0_{b} \rightarrow \Xi^0_c \pi^-$
where $\Xi^0_c \rightarrow \Xi^- \pi^+$. The mass measured for the $\Xi^0_{b}$ is 5796.7 ± 5.1(stat) ± 1.4(syst) MeV/c²,
which is consistent with our earlier result [7] but does not improve upon it. Neither of these
decay channels has been reported previously, and the reconstruction of the $\Xi^0_b$ is the first
observation of this baryon in any channel.

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FIG. 1: (a) The \( \Xi^- \pi^+ \) and (b) the \( \Xi^- \pi^+ \pi^+ \) mass distributions. The mass ranges used for the \( \Xi^0_c \) and \( \Xi^+_c \) samples are indicated by the shaded areas.

FIG. 2: (a) The \( \Xi^0_c \pi^- \) and (b) the \( \Xi^+_c \pi^- \) mass distributions. A projection of the likelihood fit is overlaid as a dashed line.