Some Recent Results on Renormalization-Group Properties of Quantum Field Theories

Robert Shrock

C. N. Yang Institute for Theoretical Physics and Department of Physics and Astronomy, Stony Brook University, Stony Brook NY 11794, USA, robert.shrock@stonybrook.edu

December 8, 2021

XXXIII International (ONLINE) Workshop on High Energy Physics
“Hard Problems of Hadron Physics: Non-Perturbative QCD & Related Quests”
November 8-12, 2021
doi:10.21468/SciPostPhysProc.?

Abstract

We discuss some higher-loop studies of renormalization-group flows and fixed points in various quantum field theories.

Contents

1 Introduction 1
2 Asymptotically Free Nonabelian Gauge Theories 2
3 RG Studies of Other Theories 5
   3.1 Finite-N Gross-Neveu Model 5
   3.2 6D $\phi^3$ Theories 5
   3.3 Studies of IR-free Theories, Including 4D U(1) and O($N$) $\lambda|\vec{\phi}|^4$ 6
4 Asymptotically Free Chiral Gauge Theories 6
5 Conclusion 7
References 7

1 Introduction

A fundamental question in quantum field theory (QFT) concerns how the running coupling of a theory changes as a function of the reference Euclidean energy/momentum scale $\mu$ where it
is measured. The variation of this coupling with \( \mu \) is described by the renormalization group (RG) beta function of the theory. Here we will discuss some results that we have obtained in this area. Much of this work was with T. A. Ryttov. We will focus mainly on vectorial asymptotically free nonabelian gauge theories in \( d = 4 \) dimensions, but also discuss some other asymptotically free theories, namely the 2D finite-\( N \) Gross-Neveu model and 6D \( \phi^4 \) theories, as well as some infrared-free theories, including U(1) gauge theory, O(N) \( \phi^4 \) theory, and chiral gauge theories.

### 2 Asymptotically Free Nonabelian Gauge Theories

Let us consider an asymptotically free (AF) vectorial nonabelian gauge theory (in \( d = 4 \) dimensions) with gauge group \( G \) and \( N_f \) massless fermions \( \psi_j, j = 1, ..., N_f \), transforming according to a representation \( R \) of \( G \). We denote the running gauge coupling as \( g(\mu) \) and define \( \alpha(\mu) \equiv g(\mu)^2/(4\pi) \) and \( a(\mu) \equiv g(\mu)^2/(16\pi^2) \). The dependence of \( \alpha(\mu) \) on \( \mu \) is described by the RG beta function, \( \beta = da(\mu)/dt \), where \( dt = d \ln \mu \). This has the series expansion

\[
\beta = -2\alpha \sum_{\ell=1}^{\infty} b_{\ell} a^\ell,
\]

where \( b_{\ell} \) is the \( \ell \)-loop coefficient. For a general operator \( O \), we denote the full scaling dimension as \( D_O \) and its free-field value as \( D_{O,\text{free}} \). The anomalous dimension of this operator, denoted \( \gamma_O \), is defined via \( D_O = D_{O,\text{free}} - \gamma_O \). The coefficients \( b_1 \) and \( b_2 \) are independent of the scheme used for regularization and renormalization and are \( b_1 = (1/3)[11C_A - 4T_f N_f] \) \([1,2]\) and \( b_2 = (1/3)[34C_A^2 - 4(5C_A + 3C_f)N_f T_f] \) \([3,4]\), where \( C_2(R) \) is the quadratic Casimir invariant, and \( T(R) \) is the trace invariant, for the representation \( R \), and we use the notation \( C_2(ad j) \equiv C_A \), \( T(R) \equiv T_f \), and \( C_2(R) \equiv C_f \). The AF condition means that \( b_1 > 0 \), i.e., \( N_f < N_u \), where \( N_u = 11C_A/(4T_f) \). Since \( a(\mu) \) is small at large \( \mu \), one can self-consistently calculate \( \beta \) as a power series in \( a(\mu) \). As \( \mu \) decreases from large values in the ultraviolet (UV) to small values in the infrared (IR), \( a(\mu) \) increases.

A situation of special interest occurs if \( \beta \) has a zero at a nonzero (physical) value \( \alpha = \alpha_{IR} \). In the asymptotically free regime, this happens if the condition \( N_u > N_f > 17C_A^2/[2(5C_A + 3C_f)T_f] \) holds, so that \( b_2 < 0 \). At the two-loop (2\( \ell \)) level, the zero in \( \beta \) occurs at \( \alpha_{IR,2\ell} = -4\pi b_1/b_2 \). If \( N_f \) is close enough to \( N_u \) that this IR zero of \( \beta \) occurs at small enough coupling so that the gauge interaction does not produce any spontaneous chiral symmetry breaking (S\( \chi \)SB), then it is an exact IR fixed point (IRFP) of the RG. The theory at this IRFP exhibits scale invariance and is inferred to exhibit conformal invariance, whence the term “conformal window” for this regime. In this IR limit, the theory is in a chirally symmetric, deconfined, nonabelian Coulomb phase (NACP). If, on the other hand, as \( \mu \) decreases and \( \alpha(\mu) \) increases toward \( \alpha_{IR} \), there is a scale \( \mu = \Lambda \) at which \( \alpha(\mu) \) exceeds a critical value, \( \alpha_{c_r} \), for the formation of a fermion condensate \( \langle \bar{\psi}\psi \rangle \) with associated S\( \chi \)SB, then the fermions gain dynamical masses of order \( \Lambda \). These fermions are then integrated out of the low-energy effective field theory operative for \( \mu < \Lambda \). In this case, \( \alpha_{IR} \) is only an approximate IRFP. We define \( N_{f,cr} \) to be the critical value of \( N_f \) such that as \( N_f \) decreases below \( N_{f,cr} \), there is S\( \chi \)SB. If \( N_f \) is only slightly less than \( N_u \), so that \( \alpha_{IR} \) is small, then the theory at the IRFP is weakly coupled and is amenable to perturbative analysis \([5]\). A case of interest for studies of physics beyond the Standard Model (BSM) is \( N_f \) slightly less than \( N_{f,cr} \). In this case, there is slow-running, quasi-conformal behavior of \( a(\mu) \) over an extended interval of \( \mu \). The dynamical breaking of the approximate scale (dilatation) symmetry then leads to a light
These results show reasonable convergence at the 4-loop and 5-loop levels, and our values of $\beta$ accord with the values $\beta_{\IR} = 0.253$, $\beta'_{\IR} = 0.295$, $\beta''_{\IR} = 0.312$, $\beta'''_{\IR} = 0.360$, and $\beta^{(4)}_{\IR} = 0.403$. As $N_f$ decreases below $N_{f,cr}$, the properties of the IR theory change qualitatively, and the perturbative calculations do not apply. A unitarity upper bound in the conformal regime is $\gamma_{\psi,\IR} < 2$ (reviewed in [27]), and studies of Schwinger-Dyson equations [28] suggest that the onset of $\chi\psi$ occurs if $\gamma_{\psi,\IR} > 1$. Thus, for a given $G$ and $R$, our higher-loop calculations of the anomalous dimension $\gamma_{\psi,\IR}$ yield estimates for $N_{f,cr}$; in turn, this information is relevant for the above-mentioned BSM theories.

There is an intensive ongoing program of research in the lattice gauge theory community to study this physics. Much work has been done for $G = SU(3)$ with $R$ equal to the fundamental representation. For this theory, $N_u = 16.5$ (where a formal continuation from physical integer $N_f$ to real $N_f$ is understood). There is not yet a consensus among lattice groups concerning the value of $N_{f,cr}$ (i.e., the lower end of the conformal window as a function of $N_f$) for this theory. As an example, we consider the case $N_f = 12$. Several lattice groups [29–34] have found that this theory is IR-conformal, while Ref. [35] has argued that it is chirally broken and hence not IR-conformal. For our 5-loop analysis, we have made use of Padé resummation methods in addition to direct analysis of series expansions. As above, we denote our $n$-loop value of $\gamma_{\psi,\IR}$ as $\gamma_{\psi,\IR,n\ell}$. We calculate $\gamma_{\psi,\IR,2\ell} = 0.773$, $\gamma_{\psi,\IR,3\ell} = 0.312$, $\gamma_{\psi,\IR,4\ell} = 0.253$, and $\gamma_{\psi,\IR,5\ell} = 0.255$. These results show reasonable convergence at the 4-loop and 5-loop levels, and our values of $\gamma_{\psi,\IR,4\ell}$ and $\gamma_{\psi,\IR,5\ell}$ are in very good agreement with the values $\gamma_{\psi,\IR} = 0.23(6)$ [33] (in accord with [31, 32]) and $\gamma_{\psi,\IR} = 0.235(46)$ [34] measured in lattice simulations. Our values are also in agreement with the range of effective values reported in [35]. For $\beta'_{\IR}$ in this $N_f = 12$ theory, as calculated via power series in the IR coupling, we find $\beta'_{2\ell} = 0.360$, $\beta'_{3\ell} = 0.295$, and $\beta'_{4\ell} = 0.282$. Again, these values show good convergence, and the 4-loop value is in very good agreement with the value $\beta'_{4\ell} = 0.26(2)$ obtained from lattice measurements [32]. In our papers we have discussed corresponding comparisons with lattice results for other gauge groups.

pseudo-Nambu-Goldstone boson, the dilaton. In a BSM application, with the Higgs boson being at least partially a dilaton, this might help to solve the fine-tuning problem of why the Higgs mass is protected against large radiative corrections.

It is of interest to investigate the properties of IRFPs in these vectorial AF gauge theories. Among these properties are the anomalous dimensions of (gauge-invariant) operators, such as $\bar{\psi}\gamma^\mu\psi \equiv \sum_{i=1}^{N_f} \bar{\psi}_i \gamma^\mu \psi_i$, denoted $\gamma_{\psi,\IR}$. In general, one can express the anomalous dimension $\gamma_{\psi,\IR}$ as the series expansion

$$\gamma_{\psi,\IR} = \sum_{\ell=1}^{\infty} c_\ell a^{\ell},$$

where $c_\ell$ is the $\ell$-loop coefficient. Evaluating this with $\alpha$ set equal to the IRFP value, calculated to a given $n$-loop $(n\ell)$ order then yields $\gamma_{\psi,\IR}$ to this order, denoted as $\gamma_{\psi,\IR,n\ell}$. Another operator of interest is $\text{Tr}(F_{\lambda\rho}F^{\lambda\rho})$, where $F_{\lambda\rho}^b$ is the field-strength tensor (with $b$ a group index). The anomalous dimension $\gamma_{\chi,\IR}$ of this operator at the IRFP satisfies $\gamma_{\chi} = -\beta'_{\IR}$, where $\beta' = d\beta / d\alpha$.

As $N_f$ decreases through the conformal regime, $d_{\IR}$ increases, motivating higher-loop calculations of anomalous dimensions. We have carried out this program of calculating the UV to IR renormalization-group evolution and anomalous dimensions at an IRFP to higher-loop order in a series of papers, many with T. A. Ryttov, including [6]- [23]. Our first calculations were at the 4-loop level [6], and subsequently, we have extended these to the 5-loop level, with inputs (in the $\overline{\text{MS}}$ scheme) up to the 5-loop level from [24, 25]. (At the 4-loop level, see also [26]). Our calculations to higher-loop order enable us to describe the IR properties of the theory throughout a larger portion of the conformal window than would be possible with the lowest-order (two-loop) results. As $N_f$ decreases below $N_{f,cr}$, the properties of the IR theory change qualitatively, and the perturbative calculations do not apply. A unitarity upper bound in the conformal regime is $\gamma_{\psi,\IR} < 2$ (reviewed in [27]), and studies of Schwinger-Dyson equations [28] suggest that the onset of $\chi\psi$ occurs if $\gamma_{\psi,\IR} > 1$. Thus, for a given $G$ and $R$, our higher-loop calculations of $\gamma_{\psi,\IR}$ yield estimates for $N_{f,cr}$; in turn, this information is relevant for the above-mentioned BSM theories.
G, fermion representations $R$, and flavor numbers $N_f$. We have also studied theories with fermions in multiple different representations [23].

Since the $b_\ell$ for $\ell \geq 3$ and the $c_\ell$ for $\ell \geq 2$ depend on the scheme used for regularization and renormalization, it is important to assess the effects of this scheme dependence. We have done this in [10–14]. This scheme dependence is a generic feature of higher-order perturbative calculations, e.g., in QCD. A scheme transformation can be expressed as a mapping between $\alpha$ and $\alpha'$, or equivalently, $a$ and $a'$, which we write as $a = a' f(a')$, where $f(a')$ is the scheme transformation function. We can write $f(a')$ as a series expansion

$$f(a') = 1 + \sum_{s=1}^{s_{\text{max}}} k_s (a')^s,$$  

where $s_{\text{max}}$ may be finite or infinite. In the new scheme, the beta function is $\beta_{a'} = -2\alpha' \sum_{\ell=1}^{\infty} b'_\ell (a')^\ell$. We have calculated the $b'_\ell$ in terms of the $b_\ell$ and $k_s$. In addition to the results $b'_1 = b_1$ and $b'_2 = b_2$, we find

$$b'_3 = b_3 + k_1 b_2 + (k_1^2 - k_2)b_1,$$  

$$b'_4 = b_4 + 2k_1 b_3 + k_1^2 b_2 + (-2k_1^3 + 4k_1k_2 - 2k_3)b_1,$$  

and so forth for higher $\ell$. We have specified a set of conditions that a physically acceptable scheme transformation must satisfy and have shown that although these can easily be satisfied in the vicinity of zero coupling, they are not automatic, and can be quite restrictive, at a nonzero coupling, as is relevant for an IRFP in an UV-free (AF) theory, or a UVFP in an IR-free theory. As part of this work, we have constructed scheme transformations that can map to a scheme with vanishing coefficients at loop level $\ell \geq 3$ in the vicinity of the origin, but we have also shown that it is more difficult to try to do this at a zero of $\beta$ away from the origin. We have applied these results to assess the degree of scheme dependence in our higher-loop calculations of anomalous dimensions at IRFPs in AF gauge theories and have shown that this dependence is small. This is similar to the experience in QCD, where calculations performed to higher order exhibited reduced scheme dependence (e.g. [36] and references therein).

The anomalous dimensions of gauge-invariant operators at the IRFP are physical and hence cannot depend on the scheme used for regularization and renormalization. However, this property is not maintained by finite-order perturbative series expansions beyond the lowest orders. It is therefore useful to calculate these anomalous dimensions in a scheme-independent (SI) manner [5,37,38]. To do this, one utilizes the fact that $\alpha_R \to 0$ as $N_f \to N_u$. Hence, one can reexpress the anomalous dimensions as series expansions in the manifestly scheme-independent variable $\Delta_f = N_u - N_f$, rather than as power series in the IR coupling:

$$\gamma_{\bar{\psi}\psi, IR} = \sum_{j=1}^{\infty} \kappa_j \Delta_f^j$$  

and

$$\beta_{IR}' = \sum_{j=1}^{\infty} d_j \Delta_f^j,$$  

where $d_1 = 0$. In general, the calculation of the coefficient $\kappa_j$ in Eq. (6) requires, as inputs, the values of the $b_\ell$ for $1 \leq \ell \leq j + 1$ and the $c_\ell$ for $1 \leq \ell \leq j$. The calculation of the coefficient $d_j$...
in Eq. (7) requires, as inputs, the values of the $b_\ell$ for $1 \leq \ell \leq j$. We denote the truncation of these series to maximal power $j = p$ as $\gamma_{\bar{\psi}\psi,IR,\Delta_f}^p$ and $\beta_\Delta^p$, respectively. With Ryttov we have calculated (i) the $\kappa_j$ up to $j = 4$, and thus the series expansion for $\gamma_{\bar{\psi}\psi,IR}$ to $O(\Delta_f^4)$, and (ii) the $d_j$ up to $j = 5$ and hence $\beta_\Delta^j$ to $O(\Delta_f^5)$ for general $G$ and $R$. We have studied a number of specific theories in detail, including the gauge groups $SU(N_c)$ with $R$ equal to the fundamental, adjoint, and rank-2 symmetric and antisymmetric tensor representations, and similarly for $SO(N_c)$ and $Sp(N_c)$ for various $N_c$. For the illustrative theory discussed above, namely $SU(3)$ with $N_f = 12$ fermions in the fundamental representation, our calculations of $\gamma_{\bar{\psi}\psi,IR}$ via Eq. (6) yield slightly larger values than our calculations via Eq. (2), and our computations of $\beta_\Delta^j$ yield slightly smaller values than those that we obtained via series expansions in the IR coupling.

An interesting feature of our scheme-independent results is that $\kappa_1$ and $\kappa_2$ are manifestly positive, and this positivity also holds for $\kappa_3$ and $\kappa_4$ for a general $G$ and all of the representations $R$ that we have studied. This leads to two monotonicity properties in the conformal regime: (i) for a fixed $p$ with $1 \leq p \leq 4$, the anomalous dimension $\gamma_{\bar{\psi}\psi,IR,\Delta_f}^p$ is a monotonically increasing function of $\Delta_f$, i.e., increases monotonically with decreasing $N_f$; (ii) for a fixed $N_f$, $\gamma_{\bar{\psi}\psi,IR,\Delta_f}^p$ is a monotonically increasing function of $p$ in the range $1 \leq p \leq 4$. From our analysis of a $\mathcal{N} = 1$ supersymmetric $SU(N_c)$ gauge theory with $N_f$ conjugate pairs of chiral superfields [19], we have found that this positivity property of the $\kappa_j$ is true for all $j$.

### 3 RG Studies of Other Theories

We have also performed higher-loop studies of RG flows and possible zeros of beta functions for other theories, including (i) the 2D finite-$N$ Gross-Neveu model [39], (ii) various $\phi^3$ theories in 6D [40,41], (iii) 4D $U(1)$ gauge theory [42], (iv) 4D nonabelian gauge theories with $N_f > N_u$ [42], and (v) 4D $O(N)$ $\lambda\bar{\phi}\phi^4$ theory [43–45]. The theories (i) and (ii) are UV-free (i.e., AF), while the theories (iii)-(v) are IR-free. In these studies, we combined direct analyses of higher-loop beta functions with Padé approximants and scheme transformations to derive results.

#### 3.1 Finite-$N$ Gross-Neveu Model

The Gross-Neveu (GN) model [46] is a 2D QFT with an $N$-component massless fermion, $\psi_j$, $j = 1,\ldots,N$ and a four-fermion interaction. This model has been of interest because it exhibits some properties similar to QCD, namely asymptotic freedom and formation of massive bound states of fermions. The model was solved exactly in the $N \to \infty$ limit in [46]. In this limit, the beta function has no IR zero. This leaves open the question of whether the beta function has an IR zero for finite $N$. We investigated this in [39], using the beta function up to the 4-loop level from [47]. We found that, where the perturbative calculation of the beta function is reliable, it does not exhibit robust evidence for an IR zero.

#### 3.2 6D $\phi^3$ Theories

$\phi^3$ theories in $d = 6$ dimensions are asymptotically free, and it is of interest to investigate whether they exhibit IRFPs. We have done this in [40] with Gracey and Ryttov, using beta functions calculated up to the 4-loop order. As before, without loss of generality, we take the matter field to be massless, since a $\phi$ field with nonzero mass $m_\phi$ would be integrated out of the low-energy...
effective theory for momentum scales $\mu < m_\phi$ and hence is not relevant for the IR limit $\mu \to 0$. We have studied $\phi^3$ theories with a real 1-component $\phi$ field and also with an $N$-component field $\phi_i$ transforming according to the fundamental representation of a global SU($N$) symmetry, with a self-interaction $\propto d_{ijk} \phi^i \phi^j \phi^k + h.c.$ For both of these theories, we find evidence against an IRFP.

An interesting study of $\phi^3$ theory in a 6D spacetime with two compact dimensions by Kisselev and Petrov is [48]. In [41], we show that if a beta function is not identically zero but has a vanishing one-loop term, then it is not, in general, possible to use scheme transformations to eliminate $\ell$-loop terms with $\ell \geq 3$ in the beta function, even in the vicinity of the origin in coupling constant space.

### 3.3 Studies of IR-free Theories, Including 4D U(1) and O($N$) $\lambda |\vec{\phi}|^4$

If the $\beta$ function of a theory is positive near zero coupling, then this theory is IR-free; as the reference scale $\mu$ decreases, the coupling decreases toward 0. As $\mu$ increases from the IR, the coupling increases, and a basic question is whether the beta function has a UV zero (in the perturbatively calculable range), which would be a UV fixed point (UVFP) of the RG.

An explicit example of a UVFP in an IR-free theory occurs in the O($N$) nonlinear $\sigma$ model in $d = 2 + \epsilon$ dimensions. From a solution of this model in the $N \to \infty$ limit, one finds, for small $\epsilon$ [49,50],

$$\beta(\xi) = \epsilon \xi \left(1 - \frac{\xi}{\xi_c}\right),$$

where $\xi$ is the effective coupling and $\xi_c = 2\pi\epsilon$. Hence, assuming that $\xi$ is small for small $\mu$, it follows that $\lim_{\mu \to \infty} \xi(\mu) = \xi_c$, so the theory has a UVFP at $\xi_c$.

Let us consider a 4D U(1) gauge theory with $N_f$ fermions with a charge $q$. This theory is IR-free, and the 1-loop and 2-loop coefficients in $\beta$ have the same sign, so there is no UV zero in $\beta$ at the maximal scheme-independent order. In [42] we investigated a possible UVFP at higher-loop order. One part of our work in [42] was an analysis of the beta function using the 5-loop coefficient [51,52]. Another part made use of exact closed-form results for $N_f \to \infty$ [53]. In [42] we also performed a corresponding investigation of possible UVFP for a nonabelian gauge theory with $N_f > N_u$. In both the U(1) and nonabelian case, we found evidence against a UVFP. Of course, in neither case does this imply that the theory has a Landau pole, because the running gauge coupling gets too large for perturbative calculations to be reliable before one actually reaches this would-be pole.

In [43–45] we investigated the RG behavior of 4D O($N$) $\lambda |\vec{\phi}|^4$ theory to six-loop order, using $b_5$ from [54] and $b_6$ from [55] (in the MS scheme). Again, for values of the interaction coupling where the perturbative (and Padé resummation) methods were applicable, we did not find robust evidence for a UVFP.

## 4 Asymptotically Free Chiral Gauge Theories

The analysis of asymptotically free chiral gauge theories is also of considerable interest. The (massless) fermion content is chosen so as to avoid any gauge anomalies, mixed gauge-gravitational anomalies, and global anomaly. As the theory flows from the UV to the IR and the coupling
grows, several possible types of behavior can occur, including (i) an exact IRFP in a conformal phase; (ii) bilinear fermion condensate formation with dynamical breaking of gauge and global symmetries; or (iii) confinement with formation of massless composite fermions. These theories have been of interest for BSM physics (e.g, [56]). Our works in this area include [57]–[62], which contain references to the extensive literature.

5 Conclusion

Studies of RG flows and possible RG fixed points in quantum field theories continue to be of great interest, both from the point of view of formal theory and for applications to BSM physics. Here we have briefly discussed some of our results on higher-loop perturbative calculations with inputs up to the five-loop level for anomalous dimensions at IR fixed points in asymptotically free nonabelian gauge theories and comparisons of these results with lattice measurements. We have also discussed our results on RG flows and investigation of possible RG fixed points for several other UV-free theories and for several IR-free theories. There are many opportunities for further work in this area.

Acknowledgements

I am grateful to T. A. Ryttov and my other coauthors for fruitful collaborations. I would also like to thank V. A. Petrov and the organizers for kindly inviting me to give this talk at the XXXIII International Workshop on High-Energy Physics at the Logunov IHEP.

Funding information The research of R.S. was supported in part by the U.S. National Science Foundation under the grants NSF-16-20628 and NSF-19-15093.

References

[1] D. J. Gross and F. Wilczek, Ultraviolet Behavior of Nonabelian Gauge Theories, Phys. Rev. Lett. 30, 1343 (1973), doi:10.1103/PhysRevLett.30.1343.

[2] H. D. Politzer, Reliable Perturbative Results for Strong Interactions?, Phys. Rev. Lett. 30, 1346 (1973), doi:10.1103/PhysRevLett.30.1346.

[3] W. E. Caswell, Asymptotic Behavior of Nonabelian Gauge Theories to Two-Loop Order, Phys. Rev. Lett. 33, 244 (1974), doi:10.1103/PhysRevLett.33.244.

[4] D. R. T. Jones, Two-Loop Diagrams in Yang-Mills Theory, Nucl. Phys. B 75, 531 (1974), doi:10.1016/0550-3213(74)90093-5.

[5] T. Banks and A. Zaks, On the Phase Structure of Vector-Like Gauge Theories with Massless Fermions, Nucl. Phys. B 196, 189 (1982), doi:10.1016/0550-3213(82)90035-9.

[6] T. A. Ryttov and R. Shrock, Higher-Loop Corrections to the Infrared Evolution of a Gauge Theory with Fermions, Phys. Rev. D 83, 056011 (2011), doi:10.1103/PhysRevD.83.056011.
[7] R. Shrock, *Higher-Loop Structural Properties of the \( \beta \) Function in Asymptotically Free Vectorial Gauge Theories*, Phys. Rev. D 87, 105005 (2013), doi:10.1103/PhysRevD.87.105005.

[8] R. Shrock, *Higher-Loop Calculations of the Ultraviolet to Infrared Evolution of a Vectorial Gauge Theory in the Limit \( N_c \to \infty, N_f \to \infty \) with \( N_f/N_c \) Fixed*, Phys. Rev. D 87, 116007 (2013), doi:10.1103/PhysRevD.87.105005.

[9] T. A. Ryttov and R. Shrock, *Infrared Zero of \( \beta \) and Value of \( \gamma_m \) for an SU(3) Gauge Theory at the Five-Loop Level*, Phys. Rev. D 94, 105015 (2016), doi:10.1103/PhysRevD.94.105015.

[10] T. A. Ryttov and R. Shrock, *An Analysis of Scheme Transformations in the Vicinity of an Infrared Fixed Point*, Phys. Rev. D 86, 085005 (2012), doi:10.1103/PhysRevD.86.085005.

[11] R. Shrock, *Study of Scheme Transformations to Remove Higher-Loop Terms in the \( \beta \) Function of a Gauge Theory*, Phys. Rev. D 88, 036003 (2013), doi:10.1103/PhysRevD.88.036003.

[12] R. Shrock, *Generalized Scheme Transformations for the Elimination of Higher-Loop Terms in the Beta Function of a Gauge Theory*, Phys. Rev. D 90, 045011 (2014), doi:10.1103/PhysRevD.90.045011.

[13] G. Choi and R. Shrock, *New Scheme Transformations and Application to Study Scheme Dependence of an Infrared Zero of the Beta Function in Gauge Theories*, Phys. Rev. D 90, 125029 (2014), doi:10.1103/PhysRevD.90.125029.

[14] G. Choi and R. Shrock, *An Integral Formalism for the Construction of Scheme Transformations in Quantum Field Theory*, Phys. Rev. D 94, 065038 (2016), doi:10.1103/PhysRevD.94.065038.

[15] T. A. Ryttov and R. Shrock, *Scheme-Independent Calculation of \( \gamma_\psi \) for an SU(3) Gauge Theory*, Phys. Rev. D 94, 105014 (2016), doi:10.1103/PhysRevD.94.105014.

[16] T. A. Ryttov and R. Shrock, *Scheme-Independent Series Expansions at an Infrared Zero of the Beta Function in Asymptotically Free Gauge Theories*, Phys. Rev. D 94 125005 (2016), doi:10.1103/PhysRevD.94.125005.

[17] T. A. Ryttov and R. Shrock, *Higher-Order Scheme-Independent Series Expansions of \( \gamma_{\psi,IR} \) and \( \beta_{IR} \) in Conformal Field Theories*, Phys. Rev. D 95, 105004 (2017), doi:10.1103/PhysRevD.95.105004.

[18] T. A. Ryttov and R. Shrock, *Infrared Fixed Point Physics in SO(\( N_c \)) and Sp(\( N_c \)) Gauge Theories*, Phys. Rev. D 96, 105015 (2017), doi:10.1103/PhysRevD.96.105015.

[19] T. A. Ryttov and R. Shrock, *Scheme-Independent Calculations of Physical Quantities in an \( \mathcal{N} = 1 \) Supersymmetric Gauge Theory*, Phys. Rev. D 96, 105018 (2017), doi:10.1103/PhysRevD.96.105018.

[20] T. A. Ryttov and R. Shrock, *Physics of the Non-Abelian Coulomb Phase: Insights from Padé Approximants*, Phys. Rev. D 97, 025004 (2018), doi:10.1103/PhysRevD.97.025004.

[21] T. A. Ryttov and R. Shrock, *Scheme-Independent Series for Anomalous Dimensions of Higher-Spin Operators at an Infrared Fixed Point in a Gauge Theory*, Phys. Rev. D 101, 076018 (2020), doi:10.1103/PhysRevD.101.076018.
[22] J. A. Gracey, T. A. Ryttov, and R. Shrock, \textit{Scheme-Independent Calculations of Anomalous Dimensions of Baryon Operators in Conformal Field Theories}, Phys. Rev. D \textbf{97}, 116018 (2018), doi:10.1103/PhysRevD.97.116018.

[23] T. A. Ryttov and R. Shrock, \textit{Scheme-Independent Series Calculations of Properties at a Conformal Infrared Fixed Point in Gauge Theories with Multiple Fermion Representations}, Phys. Rev. D \textbf{98}, 096003 (2018), doi:10.1103/PhysRevD.98.096003.

[24] P. A. Baikov, K. G. Chetyrkin, and J. H. Kühn, \textit{Five-Loop Running of the QCD Coupling Constant}, Phys. Rev. Lett. \textbf{118}, 082002 (2017), doi:10.1103/PhysRevLett.118.082002.

[25] F. Herzog, B. Ruijl, T. Ueda, J. A. M. Vermaseren, and A. Vogt, \textit{The Five-loop Beta Function of Yang-Mills Theory with Fermions}, JHEP 02(2017)090, doi:10.1007/JHEP02(2017)090.

[26] C. Pica and F. Sannino, \textit{Ultraviolet and Infrared Zeros of Gauge Theories at the Four-loop Order and Beyond}, Phys. Rev. D \textbf{83}, 035013 (2011), doi:10.1103/PhysRevD.83.035013.

[27] Y. Nakayama, \textit{Scale Invariance vs. Conformal Invariance}, Phys. Rept. \textbf{569}, 1 (2015), doi:10.1016/j.physrep.2014.12.003.

[28] T. Appelquist, K. D. Lane, and U. Mahanta, \textit{On the Ladder Approximation for Spontaneous Chiral Symmetry Breaking}, Phys. Rev. Lett. \textbf{61}, 1553 (1988), doi:10.1103/PhysRevLett.61.1553.

[29] T. Appelquist, G. Fleming, and E. Neil, \textit{Lattice Study of the Conformal Window in QCD-like Theories}, Phys. Rev. Lett. \textbf{100}, 171607 (2008), doi:10.1103/PhysRevLett.100.171607.

[30] T. Appelquist, G. Fleming, and E. Neil, \textit{Lattice Study of Conformal Behavior in SU(3) Yang-Mills Theories}, Phys. Rev. D \textbf{79}, 076010 (2009), doi:10.1103/PhysRevD.79.076010.

[31] A. Cheng, A. Hasenfratz, G. Petropoulos, and D. Schaich, \textit{Scale-Dependent Mass Anomalous Dimension from Dirac Eigenmodes}, JHEP 07(2013)061, doi:10.1007/JHEP07(2013)061.

[32] A. Hasenfratz and D. Schaich, \textit{Nonperturbative $\beta$ Function of Twelve-flavor SU(3) Gauge Theory}, JHEP 02(2018)132, doi:10.1007/JHEP02(2018)132.

[33] A. Carosso, A. Hasenfratz, and E. T. Neil, \textit{Nonperturbative Renormalization of Operators in Near-Conformal Systems Using Gradient Flows}, Phys. Rev. Lett. \textbf{121}, 201601 (2018), doi:10.1103/PhysRevLett.121.201601.

[34] M. P. Lombardo, K. Miura, T. J. Nunes da Silva, and E. Pallante, \textit{On the Particle Spectrum and the Conformal Window}, JHEP 12(2014)183, doi:10.1007/JHEP12(2014)183.

[35] Z. Fodor, K. Holland, J. Kuti, D. Nogradi, and C.-H. Wong, \textit{Fate of the Conformal Fixed Point with Twelve Massless Fermions and SU(3) Gauge Group}, Phys. Rev. D \textbf{94}, 091501 (2016), doi:10.1103/PhysRevD.94.091501.

[36] S. J. Brodsky, M. Mojaza, and X.-G. Wu, \textit{Systematic Scale-Setting to All Orders: The Principle of Maximum Conformality and Commensurate Scale Relations}, Phys. Rev. D \textbf{89}, 014027 (2014), doi:10.1103/PhysRevD.89.014027.

[37] G. Grunberg, \textit{Method of Effective Charges and Brodsky-Lepage-Mackenzie Criterion}, Phys. Rev. D \textbf{46}, 2228 (1992), doi:10.1103/PhysRevD.46.2228.
[38] T. A. Ryttov, *Consistent Perturbative Fixed Point Calculations in QCD and Supersymmetric QCD*, Phys. Rev. Lett. 117, 071601 (2016), doi:10.1103/PhysRevD.94.105014.

[39] G. Choi, T. A. Ryttov, and R. Shrock, *On the Question of a Possible Infrared Zero in the Beta Function of the Finite-N Gross-Neveu Model*, Phys. Rev. D 95, 025012 (2017), doi:10.1103/PhysRevD.95.025012.

[40] J. A. Gracey, T. A. Ryttov, and R. Shrock, *Renormalization-Group Behavior of $\phi^3$ Theories in $d = 6$ Dimensions*, Phys. Rev. D 102, 045016 (2020), doi:10.1103/PhysRevD.102.045016.

[41] T. A. Ryttov and R. Shrock, *Effect of Scheme Transformations on a Beta Function with Vanishing One-Loop Term*, Phys. Rev. D 102, 056016 (2020), doi:10.1103/PhysRevD.102.056016.

[42] R. Shrock, *Study of UV Zero of the Beta Function in Gauge Theories with Many Fermions*, Phys. Rev. D 89, 045019 (2014), doi:10.1103/PhysRevD.89.045019.

[43] R. Shrock, *On the Question of a Zero in the Beta Function of the $\lambda(\bar{\phi}^2)^2$ Theory*, Phys. Rev. D 90, 065023 (2014) doi:10.1103/PhysRevD.90.065023.

[44] R. Shrock, *Study of the Six-Loop Beta Function of the $\lambda\phi^4$ Theory*, Phys. Rev. D 94, 125026 (2016), doi:10.1103/PhysRevD.94.125026.

[45] R. Shrock, *Study of the Question of an Ultraviolet Zero in the Six-Loop Beta Function of the $O(N)$ $\lambda|\bar{\phi}|^4$ Theory*, Phys. Rev. D 96, 056010 (2017), doi:10.1103/PhysRevD.96.056010.

[46] D. J. Gross and A. Neveu, *Dynamical Symmetry Breaking in Asymptotically Free Field Theories*, Phys. Rev. D 10, 3235 (1974), doi:10.1103/PhysRevD.10.3235.

[47] J. A. Gracey, T. Luthe, and Y. Schröder, *Four Loop Renormalization of the Gross-Neveu Model*, Phys. Rev. D 94, 125028 (2016), doi:10.1103/PhysRevD.94.125028.

[48] A. V. Kisselev and V. A. Petrov, *Can Effective Four-Dimensional Scalar Theory be Asymptotically Free in a Spacetime with Higher Dimensions?*, Phys. Rev. D 103, 085012 (2021), doi:10.1103/PhysRevD.103.085012.

[49] E. Brézin and J. Zinn-Justin, *Spontaneous Breakdown of Continuous Symmetries Near Two Dimensions*, Phys. Rev. B 14, 3110 (1976), doi:10.1103/PhysRevB.14.3110.

[50] W. A. Bardeen, B. W. and R. E. Shrock, *Phase Transition in the Nonlinear $\sigma$ Model in a $2 + \epsilon$ Dimensional Continuum*, Phys. Rev. D 14, 985 (1976), doi:10.1103/PhysRevD.14.985.

[51] A. L. Kataev and S. A. Larin (2012) *Analytical Five-loop Expressions for the Renormalization Group QED $\beta$-Function in Different Renormalization Schemes*, JETP Lett. 96, 64 (2012), doi:10.1134/S0021364012130073.

[52] P. A. Baikov, K. G. Chetyrkin, J. H. Kühn, and J. Rittinger, *Vector Correlator in Massless QCD at Order $O(\alpha_s^5)$ and the QED Beta-Function at Five Loops*, JHEP 07(2012)017, doi:10.1007/JHEP07(2012)017z.

[53] A. Palanques-Mestre and P. Pascual, *The $1/N_f$ Expansion of the Gamma and Beta Functions in QED*, Commun. Math. Phys. 95, 277 (1984), doi:10.1007/BF01212398.
[54] H. Kleinert, J. Neu, V. Schulte-Frohlinde, K. G. Chetyrkin, and S. A. Larin, *Five Loop Renormalization Group Functions of O(n) Symmetric $\phi^4$ Theory and Epsilon Expansions of Critical Exponents up to $\epsilon^5*, Phys. Lett. B 272, 39 (1991), doi:10.1016/0370-2693(91)91009-K.

[55] V. Kompaniets and E. Panzer, *Minimally Subtracted Six-Loop Renormalization of O(n)-Symmetric $\phi^4$ Theory and Critical Exponents*, Phys. Rev. D 96, 036016 (2017), doi:10.1103/PhysRevD.96.036016.

[56] S. Raby, S. Dimopoulos, and L. Susskind, *Tumbling Gauge Theories*, Nucl. Phys. B 169, 373 (1980), doi:10.1016/0550-3213(80)90093-0.

[57] T. Appelquist, A. Cohen, M. Schmaltz, and R. Shrock, *New Constraints on Chiral Gauge Theories*, Phys. Lett. B 459, 235 (1999), doi:10.1016/S0370-2693(99)00616-4.

[58] T. Appelquist and R. Shrock, *Neutrino Masses in Theories with Dynamical Electroweak Symmetry Breaking*, Phys. Lett. B 548, 204 (2002), doi:10.1016/S0370-2693(02)02854-X.

[59] T. Appelquist and R. Shrock, *Dynamical Symmetry Breaking of Extended Gauge Symmetries*, Phys. Rev. Lett. 90, 201801 (2003), doi:10.1103/PhysRevLett.90.201801.

[60] T. Appelquist and R. Shrock, *On the Ultraviolet to Infrared Evolution of Chiral Gauge Theories*, Phys. Rev. D 88, 105012 (2013), doi:10.1103/PhysRevD.88.105012.

[61] Y. Shi and R. Shrock, *Renormalization-Group Evolution and Nonperturbative Behavior of Chiral Gauge Theories with Fermions in Higher-Dimensional Representations*, Phys. Rev. D 92, 125009 (2015), doi:10.1103/PhysRevD.92.125009.

[62] Y. Shi and R. Shrock, *$A_k \hat{F}$ Chiral Gauge Theories*, Phys. Rev. D 92, 105032 (2015), doi:10.1103/PhysRevD.92.105032.