Controlling the Kondo Effect in CoCu\textsubscript{n} Clusters Atom by Atom

N. Néel,\textsuperscript{1} J. Kröger,\textsuperscript{1} R. Berndt,\textsuperscript{1} T. Wehling,\textsuperscript{2} A. Lichtenstein,\textsuperscript{2} and M. I. Katsnelson\textsuperscript{3}

\textsuperscript{1}Institut für Experimentelle und Angewandte Physik, Christian-Albrechts-Universität zu Kiel, D-24098 Kiel, Germany
\textsuperscript{2}Institut für Theoretische Physik I, Universität Hamburg, D-20355 Hamburg, Germany
\textsuperscript{3}Institute for Molecules and Materials, Radboud University Nijmegen, NL-6525 AJ Nijmegen, The Netherlands

(Dated: October 1, 2008)

Clusters containing a single magnetic impurity were investigated by scanning tunneling microscopy, spectroscopy, and \textit{ab initio} electronic structure calculations. The Kondo temperature of a Co atom embedded in Cu clusters on Cu(111) exhibits a non-monotonic variation with the cluster size. Calculations model the experimental observations and demonstrate the importance of the local and anisotropic electronic structure for correlation effects in small clusters.

PACS numbers: 68.37.Ef,72.15.Qm,73.20.Fz

Nanometer scaled electronic devices require the understanding and the control of electron behavior, in particular of correlation effects, at the atomic scale. Experimental techniques such as scanning tunneling microscopy (STM) and spectroscopy and angle-resolved photoemission demonstrate the relevance of many-body phenomena beyond standard band theory \cite{1}. In many cases, especially for compounds of \textit{d} and \textit{f} elements, strong electron correlations should be taken into account even for a qualitative description of electron energy spectra and physical properties \cite{2}. The Kondo effect, which is related to the resonant scattering of conduction electrons by quantum local centers, is one of the key correlation effects in condensed matter physics \cite{3}. It arises from the interaction of a single localized magnetic moment with a continuum of conduction electrons and results in a sharp resonance likewise serves as an indicator of the cluster state whose energy decreases with increasing number of state whose energy decreases with increasing number of conduction electron states at the magnetic site.

Experiments were performed with a home-made scanning tunneling microscope operated at 7 K and in ultrahigh vacuum with a base pressure of 10\textsuperscript{-9} Pa. Tungsten tips and Cu(111) surfaces were prepared by argon ion bombardment and annealing. While single Co atoms were deposited onto the cold surface using an electron beam evaporator and an evaporant of 99.99\% purity, single Cu atoms were transferred from the tip as previously reported \cite{17}. The individual atoms were chemically identified by the presence (Co) or absence (Cu) of the Abrikosov-Suhl resonance. Clusters consisting of a single Co atom and several Cu atoms \cite{13,18} were fabricated by atom manipulation with the microscope tip. Spectroscopy was performed by a state-of-the-art lock-in technique.

Figures\textsuperscript{4}(a)-(e) show a series of constant-current STM images, which illustrate fabrication of a CoCu cluster [Figs\textsuperscript{4}(a), (b)] and dimensions of CoCu\textsubscript{n} (\textit{n} = 1, \ldots, 4) clusters [Figs\textsuperscript{4}(b)-(e)]. In STM images acquired at a sample voltage of 100 mV Co atoms appear higher than Cu atoms [Fig\textsuperscript{4}(a)], and thus an additional means of discriminating the adatom species is provided. The stick-and-ball models of the clusters used for the calculations are presented in Fig\textsuperscript{4}(f). Spectroscopy of the differential conductance \((dI/dV)\) performed with the tip positioned above the cluster center reveals an unoccupied state whose energy decreases with increasing number of Co atoms [squares in Fig\textsuperscript{4}(g)]. Our calculations show that this resonance is of \textit{p}\textsubscript{z} character and are in agreement with experimental data [circles in Fig\textsuperscript{4}(g)]. This resonance likewise serves as an indicator of the cluster
FIG. 1: (Color online) Scanning tunneling microscopy images and resonance energies of CoCu$_n$ clusters acquired at 7 K. (a) STM images of single Co and Cu atoms adsorbed on Cu(111) prior to fabricating CoCu$_n$ clusters. (b) - (e) Sequence of STM images showing CoCu$_n$ ($n = 1, \ldots, 4$) clusters. The same lateral and height scale was used for each image. Images were acquired at a sample voltage of 100 mV and a tunneling current of 0.1 nA. (f) Stick-and-ball models of fully relaxed adsorption structures of CoCu$_n$ ($n = 0, \ldots, 4$) on Cu(111). (g) Experimental energy of unoccupied cluster state versus the cluster size. The size of the symbols (squares) corresponds to the uncertainty of the energies. Calculated energies are depicted as circles.

Figure 2 shows a sequence of $dI/dV$ spectra acquired on a single Co atom (lower curve, $n = 0$) and above the center of CoCu$_n$ clusters with $n$ ranging between 1 and 4. Starting from a sharp indentation of the differential conductance close to the Fermi level (sample voltage $V = 0$ mV) which is the spectroscopic signature of the single-Co Kondo effect on Cu(111), the $dI/dV$ curve broadens appreciably upon adding two Cu atoms to the Co atom. Surprisingly, upon increasing the number of hybridizations to three and four Cu atoms, $dI/dV$ spectra exhibit a sharpening of the resonance again in contrast with the monotonic behaviour of the unoccupied $p_z$-like state with the number of copper atoms [Fig. 1(g)].

To quantify the broadening of the resonances at the Fermi level we described experimental data by a Fano line shape \[ 1 \]:

$$
\frac{dI}{dV} \propto \frac{(q + \epsilon)^2}{1 + \epsilon^2}
$$

with $q$ the asymmetry parameter of the Fano theory and $\epsilon = (eV - \epsilon_F)/(k_B T_K)$ where $-e$ denotes the electron charge, $V$ the sample voltage, $\epsilon_F$ the resonance energy, and $k_B$ Boltzmann’s constant. The profiles of Fano line shapes according to Eq. 1 and fitted to $dI/dV$ data are presented as solid lines in Fig. 2. The striking behavior of broadening and sharpening of the Abrikosov-Suhl resonance with increasing number of Cu atoms is reflected by the Kondo temperature $T_K$ which we plotted in Fig. 3(a). While the single Co adatom exhibits $T_K = (61 \pm 4)$ K, which is in agreement with a previous publication \[3\], for CoCu$_2$ we find $(326 \pm 30)$ K which then decreases again to $(43 \pm 6)$ K for CoCu$_4$. Figure 3(b) shows the evolution of the asymmetry parameter with the cluster size. For $n = 0, 1, 2$ the asymmetry parameter varies weakly around 0.1. It increases steeply by $\approx 100 \%$ upon adding a third Cu atom. Kondo temperatures and asymmetry parameters for the various clusters are summarized in Table I.

Table I: Experimental and calculated Kondo temperatures ($T_K$) and asymmetry parameters ($q$) for clusters CoCu$_n$. $T_K$ and $q$ were extracted from fits of the Fano line shape (Eq. 1) to experimental data. Prior to spectroscopy the tunneling gap was set at 30 mV and 1 nA. Spectra for $n = 1, 2, 3, 4$ are vertically offset by 5, 10, 20, 25 nS, respectively.

| $n$ | $T_K$ (K) | $q$       | $T_K$ (K) | $q$       |
|-----|-----------|-----------|-----------|-----------|
| 0   | 61 ± 4    | 0.09 ± 0.03 | 52 ± 4    | 0.12 ± 0.01 |
| 1   | 142 ± 10  | 0.12 ± 0.03 | 204 ± 12  | 0.03 ± 0.01 |
| 2   | 326 ± 30  | 0.07 ± 0.03 | 288 ± 3   | 0.01 ± 0.01 |
| 3   | 200 ± 15  | 0.22 ± 0.04 | 221 ± 18  | 0.13 ± 0.01 |
| 4   | 43 ± 6    | 0.22 ± 0.04 | 79 ± 3    | 0.33 ± 0.01 |

The peculiar behavior of the Kondo resonance with cluster size is at variance with the monotonic dependence...
FIG. 3: (Color online) Experimental (squares) and calculated (circles) Kondo temperatures and asymmetry parameters. (a) Kondo temperature \( T_K \) as a function of cluster size. (b) Asymmetry parameter \( q \) as a function of cluster size.

The Kondo temperature for a single magnetic impurity can be estimated as

\[
T_K \approx W \sqrt{|J N(E_F)|} \exp \left( -\frac{1}{|J N(E_F)|} \right),
\]

where \( W \) is the conduction \((s)\) electron bandwidth, \( J \) the \( s-d \) exchange integral, and \( N(E_F) \) the density of states at the Fermi energy \( E_F \). Since the \( s-d \) exchange interaction is local, \( N(E_F) \) in Eq. 2 is the local density of states of conduction electrons at the magnetic atom.

The asymmetry parameter, \( q \), is also determined by the local electronic structure. In a simple model \( q \) may be expressed as

\[
q = \frac{\gamma + \text{Re} G(E_F)}{\text{Im} G(E_F)},
\]

where \( G \) is the local Green’s function of the conduction electrons at the impurity site and \( \gamma \) measures the ratio of the coupling of the scanning tunneling microscope tip to conduction electron states and to the strongly localized Co \( d \) states.

To realistically describe the CoCuₙ clusters on Cu(111) we calculated their electronic and structural properties by means of density functional theory using the generalized gradient approximation to the exchange correlation potential. For solving the resulting Kohn-Sham equations accurately, the Vienna Ab Initio Simulation Package with the projector augmented wave basis sets and 350 eV as plane wave cut-off have been used. We modeled the CoCuₙ structures using \( 4 \times 3 \) supercells of Cu(111) slabs containing up to 7 Cu layers. These structures have been fully relaxed with the requirement that all forces are less than 0.2 eV nm\(^{-1} \) and
were then used to calculate the local density of states. The supercell Brillouin zone integrations for obtaining the local density of states have been performed using the tetrahedron-method on $6 \times 6 k$ meshes with subsequent 50 - 100 meV Gaussian broadening. Below, we use both, 50 and 100 meV smearing, to provide an uncertainty margin for the derived Kondo temperatures and asymmetry parameters.

The relaxed structures [Fig. 1(f)] show that the distance of the Co atom to the Cu surface increases monotonically upon adding Cu atoms to the cluster. However, the local electronic structure at the Fermi level and at the Co site varies in a non-monotonic way with increasing number of Cu atoms. The substrate conduction band of Co site varies in a non-monotonic way with increasing number of Cu atoms. The substrate conduction band states extend to the Co adatom site where they interact with localized Co 3d electrons. This interaction is quantified by the Co $sp$-projected local density of states, $N(E)$, which is shown in Figs. 1(a) and 1(b). The dependence of $G(E_F)$ on the number of Cu atoms is non-monotonic [see the close-up view around $E_F$ in Fig. 1(b)] and leads, according to Eq. 2, to the experimentally observed non-monotonic behavior of the Kondo temperature (Fig. 3).

We used $W = 20 \text{ eV}$ and $J = 1.3 \text{ eV}$ as fitting parameters. Moreover, with $\gamma = 0.22 \text{ eV}^{-1}$ in Eq. 3 the variation of $G(E_F)$ with the number Cu atoms also reproduces the non-monotonic trend of $\gamma$ observed in the experiments [Fig. 3(b)]. Figure 3(c) shows the evolution of unoccupied resonance energies with the number of Cu atoms. The resonances appear as peaks in the calculated $p_z$-projected local density of states whose energies show the same monotonic behavior as in the experiment [Fig. 3(g)].

The non-monotonic dependence of $N(E_F)$ on the number of Cu atoms surrounding the Co atom is related to the specific nature of chemical bonds in the system. Since localized $d$ orbitals of Co atoms are strongly anisotropic and form well-defined directional bonds the degree of their hybridization with Cu atoms is determined by the whole geometry of the cluster rather than by just the coordination number.

In summary, we have shown that the Kondo effect in clusters is but depends crucially on their detailed geometric structure. Atom by atom manipulation changes the local density of conduction electron states at the magnetic site and thus varies the Kondo temperature in rather broad limits. The results demonstrate that electron correlations may be tuned in atomic-scale structures.

Discussion with H. Kroha (University of Bonn, Germany) and financial support by the Deutsche Forschungsgemeinschaft through SFB 668 and SPP 1153 are gratefully acknowledged.

* Electronic address: kroeger@physik.uni-kiel.de

[1] E. W. Plummer, Ismail, R. Matzdorf, A. V. Melechko, J. P. Pierce, and J. Zhang, Surf. Sci. 500, 1 (2002).
[2] G. Koi, S. Y. Savrasov, K. Haule, V. S. Oudovenko, O. Parcollet, and C. A. Marianetti, Rev. Mod. Phys. 78, 865 (2006).
[3] A. C. Hewson, The Kondo problem to heavy fermions, (Cambridge University Press, 1993).
[4] D. Goldhaber-Gordon, H. Shtrikman, D. Mahalu, D. Abusch-Magder, U. Meirav, and M. A. Kastner, Nature 391, 156 (1998).
[5] S. M. Crone, T. H. Oosterkamp, and L. P. Kouwenhoven, Science 309, 540 (1998).
[6] J. Li, W.-D. Schneider, and R. Berndt, Phys. Rev. Lett. 80, 2893 (1998).
[7] V. Madhavan, W. Chen, T. Jamneala, M. F. Crommie, and N. S. Wingreen, Science 280, 567 (1998).
[8] H. C. Manoharan, C. P. Lutz, and D. M. Eigler, Nature 403, 512 (2000).
[9] N. Knorr, M. A. Schneider, L. Diekhöner, P. Wahl, and K. Kern, Phys. Rev. Lett. 88, 096804 (2002).
[10] A. Zhao, Q. Li, L. Chen, H. Xiang, W. Wang, S. Pan, B. Wang, X. Xiao, J. Yang, J. G. Hou, and Q. Zhu, Science 309, 1542 (2005).
[11] J. Kliever, R. Berndt, and S. Crampin, Phys. Rev. Lett. 85, 4936 (2000).
[12] N. Nilius, T. M. Wallis, and W. Ho, Science 297, 1853 (2002).
[13] J. Lagoute, C. Nacci, and S. Fölsch, Phys. Rev. Lett. 98, 146804 (2007).
[14] A. J. Heinrich, J. A. Gupta, C. P. Lutz, and D. M. Eigler, Science 306, 466 (2004).
[15] T. Jamneala, V. Madhavan, and M. F. Crommie, Phys. Rev. Lett. 87, 256804 (2001).
[16] U. Heiz, A. Sanchez, S. Abbet, and W.-D. Schneider, W.-D., J. Am. Chem. Soc. 121, 3214 (1999).
[17] L. Limot, J. Kröger, R. Berndt, A. García-Lekue, and W. A. Hofer, Phys. Rev. Lett. 94, 26102 (2005).
[18] A. J. Stroscio, F. Tavazza, J. N. Crain, R. J. Celotta, and A. M. Chaka, Science 313, 948 (2006).
[19] U. Fano, Phys. Rev. 124, 1866 (1961).
[20] O. Úságyh, J. Kroha, L. Szunyogh, and A. Zawadowski, Phys. Rev. Lett. 85, 2557 (2000).
[21] C.-Y. Li, A. H. Castro Neto, and B. A. Jones, Phys. Rev. Lett. 97, 156102 (2006).
[22] G. A. Fiete and E. J. Heller, Rev. Mod. Phys. 75, 933 (2003).
[23] O. Agam and A. Schiller, Phys. Rev. Lett. 86, 484 (2001).
[24] D. Porrás, J. Fernandez-Rossier, and C. Tejedor, Phys. Rev. B 63, 155406 (2001).
[25] C. Baldacchini, L. Chiodo, F. Allegretti, C. Mariani, M. G. Betti, P. Monachesi, and R. Del Sole, Phys. Rev. B 68, 195109 (2003).
[26] L. Limot and R. Berndt, Appl. Surf. Sci. 237, 572 (2004).
[27] V. Madhavan, W. Chen, T. Jamneala, M. F. Crommie, and N. S. Wingreen, Phys. Rev. B 64, 165412 (2001).
[28] J. P. Perdew, J. A. Chevary, S. H. Vosko, K. A. Jackson, M. R. Pederson, D. J. Singh, and C. Fiolhais, Phys. Rev. B 46, 6671 (1992).
[29] G. Kresse and J. Hafner, J. Phys.: Condens. Matter 6, 8245 (1994).
[30] G. Kresse and D. Joubert, Phys. Rev. B 59, 1758 (1999).
[31] P. E. Blöchl, Phys. Rev. B 50, 17953 (1994).