COMPUTING HODGE INTEGRALS WITH ONE $\lambda$-CLASS

YON-SEO KIM

Department of Mathematics, UCLA

Abstract. All Hodge integrals with at-most one $\lambda$-class can be expressed as polynomials in terms of lower dimensional Hodge integrals with at-most one $\lambda$-class. Algorithm to compute any given Hodge integral with at-most one $\lambda$-class is discussed and some examples are presented.

1. Introduction

Let $\overline{M}_{g,n}$ denote the Deligne-Mumford moduli stack of stable curves of genus $g$ with $n$ marked points. A Hodge integral is an integral of the form

$$\int_{\overline{M}_{g,n}} \psi_1^{j_1} \ldots \psi_n^{j_n} \lambda_1^{k_1} \ldots \lambda_g^{k_g}$$

where $\psi_i$ is the first Chern class of the cotangent line bundle at the $i$-th marked point, and $\lambda_1, \ldots, \lambda_g$ are the Chern classes of the Hodge bundle. Hodge integrals arise naturally in the calculations of Gromov-Witten invariants by localization techniques. Their explicit evaluations are difficult problems. The famous Witten's conjecture/Kontsevich’s theorem [21], [10] gives a recursive relation of Hodge integrals involving $\psi$ classes only:

$$\int_{\overline{M}_{g,n}} \psi_1^{j_1} \ldots \psi_n^{j_n}$$

and some of them can be computed recursively through String Equation and KdV hierarchy. In [4], C.Faber developed an algorithm to compute intersection numbers of type [1]. Also, E.Getzler obtained recursion relations [3] for the case of $g = 2$, one of which is given as:

$$\langle \langle \tau_{k+2}(x) \rangle \rangle_2 = \langle \langle \tau_{k+1}(x) \gamma_a \rangle \rangle_0 \langle \langle \gamma^a \rangle \rangle_2 + \langle \langle \tau_k(x) \gamma_a \rangle \rangle_0 \langle \langle \gamma^b \rangle \rangle_2 + \frac{7}{10} \langle \langle \tau_k(x) \gamma_a \gamma_b \rangle \rangle_0 \langle \langle \gamma^a \rangle \rangle_1 \langle \langle \gamma^b \rangle \rangle_1$$

$$+ \frac{1}{10} \langle \langle \tau_k(x) \gamma_a \gamma_b \rangle \rangle_0 \langle \langle \gamma^a \gamma^b \rangle \rangle_1 - \frac{1}{240} \langle \langle \tau_k(x) \gamma_a \rangle \rangle_1 \langle \langle \gamma^a \gamma_b \gamma^b \rangle \rangle_0$$

$$+ \frac{13}{240} \langle \langle \tau_k(x) \gamma_a \gamma^a \gamma_b \rangle \rangle_0 \langle \langle \gamma^b \rangle \rangle_1 + \frac{1}{960} \langle \langle \tau_k(x) \gamma_a \gamma^a \gamma_b \gamma^b \rangle \rangle_0$$

for $k \geq 0$

When the $\lambda$-classes are involved, the computation of Hodge integrals is not easy. There was $\lambda_g$-conjecture which computes for the case of one top-degree $\lambda$-class as [4];
\( \int_{\overline{M}_{g,n}} \psi_1^{k_1} \cdots \psi_n^{k_n} \lambda_y = \left( \frac{2g + n - 3}{k_1, \ldots, k_n} \right) \frac{2^{2g-1} - 1}{2^{2g-1}} \frac{|B_{2g}|}{(2g)!} \)

where \( B_{2g} \) are Bernoulli numbers and \( k_1 + \cdots + k_n = 2g - 3 + n \). By using Mariño-Vafa formula \([13][14]\), the case of \( \lambda_y = 1 \) with one marked point can be computed as;

\[ \int_{\overline{M}_{g,1}} \psi_1^{2g-1} \lambda_y - 1 = b_g \sum_{i=1}^{2g-1} \frac{1}{i} - \frac{1}{2} \sum_{g_1 + g_2 = g} \frac{(2g_1 - 1)!(2g_2 - 1)!}{(2g - 1)!} b_{g_1} b_{g_2} \]

and the case of more than one marked points can be computed by repeatedly applying the Cut-and-Join equation:

\[ \frac{\partial \Omega}{\partial \tau} = \frac{\sqrt{-1}}{2} \sum_{i,j \geq 1} (ijp_i + j \partial^2 \Omega / \partial p_i \partial p_j + ijp_i + j \partial \Omega / \partial p_i \partial p_j + (i + j)p_ip_j \partial \Omega / \partial p_{i+j}) \]

which was used in the proof of Mariño-Vafa formula, and proven to be an effective tool in studying Hodge integrals.

The moduli space of relative stable morphisms admits a natural \( S^1 \)-action induced from the \( S^1 \)-action on the target space. And as a result of the localization formula applied to it, the following convolution formula is obtained;

**Theorem 6.1.** For any partition \( \mu \) and \( e \) with \( |e| < |\mu| + l(\mu) - \chi \), we have

\[ [\lambda^{(\mu)}] \sum_{|\nu|=|\mu|} \Phi_{\nu,\nu}^\lambda (-\lambda) z_{\nu} D_{\nu,e}^\lambda = 0 \]

where the sum is taken over all partitions \( \nu \) of the same size as \( \mu \).

Here \( \chi \) is the prescribed Euler number of domain curves, \([\lambda^a]\) means taking the coefficient of \( \lambda^a \), \( \Phi^\lambda(\lambda) \) is a generating series of Double Hurwitz Numbers, and \( D^\lambda(\lambda) \) is a certain generating series of Hodge integrals. This formula gives many relations between Hodge integrals with at-most one \( \lambda \)-class, and it is enough to consider the special case of \( \mu = (d) \) for some positive integers \( d \) to compute all Hodge integrals with at-most one \( \lambda \)-class. More precisely, the following theorem is proved in Section 7;

**Theorem 7.2.** Any given Hodge integral with one \( \lambda \)-class:

\[ \int_{\overline{M}_{g,n}} \psi_1^{k_1} \cdots \psi_n^{k_n} \lambda_j \]

where \( k_1, \ldots, k_n \in \mathbb{N} \cup \{0\}, j \in \{0, 1, 2, \ldots, g\} \), is explicitly expressed as a polynomial in terms of lower-dimensional Hodge integrals with one \( \lambda \)-class. Therefore it computes all Hodge integrals with one \( \lambda \)-class.

The rest of the paper is organized as follows: In Section 2, we summarize various versions of localization formulas which will be used in computing Hodge integrals. In Section 3, the Relative Moduli Space is defined and the natural \( S^1 \)-action on it is
introduced. Also the fixed locus of the $S^1$-action and their corresponding description in terms of graphs is discussed. In Section 4, we compute the Euler class of the normal bundle of the fixed locus of the $S^1$-action in the relative moduli space. In Section 5, the Double Hurwitz Numbers and its description in terms of Hodge integrals over a certain moduli space is discussed. In Section 6, the Recursion Formula which gives relations between Hodge integrals with at-most one $\lambda$-class is proved. In Section 7, the Recursion Formula is used prove that all Hodge integrals with at-most one $\lambda$-class is explicitly expressed as a polynomial in terms of lower dimensional Hodge integrals with at-most one $\lambda$-class. In Section 8, an algorithm to implement the Recursion Formula and to compute each Hodge integral with at-most one $\lambda$-class is discussed. In Section 9, some examples of the algorithm in section 8 are presented.

Acknowledgements. I would like to thank my advisor, Professor Kefeng Liu for his constant support and encouragements as well as many inspiring discussions. I would also like to thank Xiaowei Wang for helpful discussions and friendship.

2. Localization Formula

In this section, I will summarize various versions of localization formulas.

2.1. Equivariant Cohomology. Let $G$ be a compact Lie group acting on $M$. The equivariant cohomology of $M$ is defined as the ordinary cohomology of the space $M_G$ obtained from a fixed universal $G$-bundle $EG$, by the mixing construction

$$M_G = EG \times_G M$$

Here, $G$ acts on the right of $EG$ and on the left of $M$, and the notation means that we identify $(pg, q) \sim (p, gq)$ for $p \in EG$, $q \in M$, $g \in G$. Hence $M_G$ is the bundle with fibre $M$ over the classifying space $BG$ associated to the universal bundle $EG \longrightarrow BG$.

We have natural projection map $\pi : M_G \longrightarrow BG$ and $\sigma : M_G \longrightarrow M/G$, which fits into the mixing diagram of Cartan and Borel:

$$
\begin{array}{cccc}
E G & \longrightarrow & E G \times M & \longrightarrow & M \\
\downarrow & & \downarrow & & \downarrow \\
B G & \longrightarrow & \pi E \times_G M & \sigma & M / G
\end{array}
$$

If $G$ acts smoothly on $M$, then we have $M_G \cong M/G$. This is not true in general but it turns out that $M_G$ is a better functorial construction and the proper homotopy theoretic quotient of $M$ by $G$. In any case, the equivariant cohomology, denoted by $H^*_G(M)$, is defined by

$$H^*_G(M) = H^*(M_G)$$

and constitutes a contravariant functor from $G$-spaces to modules over the base ring $H^*_G := H^*_G(pt) = H^*(BG)$. The map $\sigma$ defines a natural map $\sigma^* : H^*(M/G) \longrightarrow H^*_G(M)$ which is an isomorphism if $G$ acts freely. The inclusion $i : M \longrightarrow M_G$ induces a natural map $i^* : H^*_G(M) \longrightarrow H^*(M)$.
2.2. Atiyah-Bott Localization Formula. Let \( i : V \hookrightarrow M \) be a map of compact manifolds. The tubular neighborhood of \( V \) inside \( M \) can be identified with the normal bundle of \( V \). On the total space of the normal bundle, there is the Thom form \( \Phi_V \) which has compact support in the fibres and integrates to one in each fiber. Extending this form by zero gives a form in \( M \), and multiplying by \( \Phi_V \) provides a map \( H^*(V) \cong H^{*+k}(M, M\setminus V) \rightarrow H^*(M) \). In particular, the cohomology class \( 1 \in H^0(V) \) is sent to the Thom class and this class restricts to be the Euler class of the normal bundle of \( V \) in \( M, N_V/M \). Hence, we see that

\[
i^*i_*1 = c(N_V/M)
\]

This also holds in equivariant cohomology by same argument applied to \( V_G, M_G \). The theorem of Atiyah and Bott says that an inverse of the Euler class always exists along the fixed locus of a group action. Precisely, \( i^*/e(N_V/M) \) is the inverse of \( i_* \) in equivariant cohomology, i.e. for any equivariant class \( \phi \),

\[
\phi = \sum_F i_* i^* \phi \quad e(T(F/X))
\]

holds where \( F \) runs over the fixed locus of the group action. In the integrated form, we have

\[
\int_M \phi = \sum_F \int_F i^* \phi \quad e(N_F/M)
\]

2.3. Functorial Localization Formula. Let \( X \) and \( Y \) be \( T \)-manifolds. Assume \( f : X \rightarrow Y \) is a \( T \)-equivariant map, \( j_E : E \hookrightarrow Y \) is a fixed component in \( Y \), and \( i_F : F \hookrightarrow f^{-1}(E) \) is a fixed component in \( X \). For any equivariant class \( \omega \in H^*_T(X) \), we have the diagrams:

\[
\begin{array}{ccc}
F & \xrightarrow{i_F} & X \\
\downarrow{g = f|_F} & & \downarrow{f} \\
E & \xrightarrow{j_E} & Y
\end{array}
\quad \quad \quad 
\begin{array}{ccc}
i^*_F(\omega) & \xrightarrow{\tilde{e}_T(F/X)} & i^*_F \\
\downarrow{g_!} & & \downarrow{f_!} \\
\tilde{e}_T[F/(X)] & \xrightarrow{i^*_E(\omega)} & \tilde{e}_T(E/Y)
\end{array}
\]

Applying the Atiyah-Bott Localization Formula with the naturality relation \( f_!(\omega \cdot f^*\alpha) = f_!\omega \cdot \alpha \), we obtain the Functorial Localization Formula:

\[
(5) \quad g_! \left[ \frac{i^*_F(\omega)}{\tilde{e}_T(F/X)} \right] = \frac{j^*_E f_!(\omega)}{\tilde{e}_T(E/Y)}
\]
2.4. Virtual Functorial Localization Formula. The above Functorial Localization Formula is also valid in the case where $X$ and $F$ are virtual fundamental classes. In this paper, we will use $\overline{\mathcal{M}}_{X,n}(\mathbb{P}^1, \mu)^{\text{vir}}$ for $X$, and $F_T^{\text{vir}}$ for $F$. Hence for any equivariant class $\omega$, we have:

\[ \int_{\overline{\mathcal{M}}_{X,n}(\mathbb{P}^1, \mu)^{\text{vir}}} \omega = \sum_{F_T} \int_{F_T^{\text{vir}}} \frac{i^*_{\text{vir}}(\omega)}{e_T(F_T/\overline{\mathcal{M}}_{X}(\mathbb{P}^1, \mu))} \]

3. Relative Moduli Space and $S^1$-action

3.1. Moduli space of relative morphisms. For any non-negative integer $m$, let

\[ \mathbb{P}^1[m] = \mathbb{P}^1(0) \cup \mathbb{P}^1(1) \cup \cdots \cup \mathbb{P}^1(m) \]

be a chain of $m + 1$ copies of $\mathbb{P}^1$ such that $\mathbb{P}^1(l)$ is glued to $\mathbb{P}^1(l+1)$ at $p_{1}^{(l)}$ for $0 \leq l \leq m - 1$. The irreducible component $\mathbb{P}^1(0)$ is referred to as the root component, and other irreducible components are called bubble components. Two points $p_{1}^{(l)} \neq p_{1}^{(l+1)}$ in $\mathbb{P}^1(l)$ are fixed. Denote by $\pi[m] : \mathbb{P}^1[m] \rightarrow \mathbb{P}^1$ the map which is identity on the root component and contracts all bubble components to $\mathbb{P}^1(0)$. Also denote by $\mathbb{P}^1(m) = \mathbb{P}^1(1) \cup \cdots \cup \mathbb{P}^1(m)$ the union of bubble components of $\mathbb{P}^1[m]$. For a fixed partition $\mu$ of a positive integer $d$, let $\overline{\mathcal{M}}_{X,n}(\mathbb{P}^1, \mu)$ be the moduli space of relative morphisms $f : (C; x_1, \cdots, x_{l(\mu)}, z_1, \cdots, z_n) \rightarrow (\mathbb{P}^1[m], p_{1}^{(m)})$ such that

1. $(C; x_1, \cdots, x_{l(\mu)}, z_1, \cdots, z_n)$ is a possibly-disconnected prestable curve of Euler number $\chi$ with $l(\mu) + n$ marked points. Here, the marked points are unordered.
2. $f^{-1}(p_{1}^{(m)}) = \sum_{i=1}^{l(\mu)} \mu_i \cdot x_i$ as Cartier divisors and $\deg(\pi[m] \circ f) = |\mu|$.
3. The preimage of each node in $\mathbb{P}^1[m]$ consists of nodes of $C$. If $f(y) = p_{1}^{(l)}$ and $C_1$ and $C_2$ are two irreducible components of $C$ which intersects at $y$, then $f |_{C_1}$ and $f |_{C_2}$ have the same contact order to $p_{1}^{(l)}$ at $y$.
4. The automorphism group of $f$ is finite. Here, an automorphism of $f$ consists of an automorphism of the domain curve and an automorphism of the pointed curve $(\mathbb{P}^1(m), p_{1}^{(0)}, p_{1}^{(m)})$.

Following [11][12], $\overline{\mathcal{M}}_{X,n}(\mathbb{P}^1, \mu)$ is a separated, proper Deligne-Mumford stack with a perfect obstruction theory of virtual dimension $r = -\chi + |\mu| + l(\mu) + n$, and hence has a virtual fundamental class of degree $r$.

3.2. Torus Action. Consider the $\mathbb{C}^*$-action $t \cdot [z^0 : z^1] = [tz^0 : z^1]$ on $\mathbb{P}^1$. There are two fixed points $p_0 = [0 : 1]$ and $p_1 = [1 : 0]$. Extend this action to the action on $\mathbb{P}^1[m]$ by identifying the root component with $\mathbb{P}^1$ and giving trivial actions on bubble components. Then this extended action on $\mathbb{P}^1[m]$ induces an action on $\overline{\mathcal{M}}_{X,n}(\mathbb{P}^1, \mu)$.
3.3. Fixed Locus. The connected components of the fixed points set of $\overline{\mathcal{M}}_{g,n}^\chi(\mathbb{P}^1, \mu)$ under the induced torus action can be parametrized by labeled graphs. For any $f : (C; x_1, \ldots, x_{\ell(\mu)}, z_1, \ldots, z_n) \to \mathbb{P}^1[m]$ representing a fixed point of the $\mathbb{C}^*$-action on $\overline{\mathcal{M}}_{g,n}^\chi(\mathbb{P}^1, \mu)$, the restriction of $\hat{f} := \pi[m] \circ f : C \to \mathbb{P}^1$ to an irreducible component of $C$ is either a constant map to one of the $\mathbb{C}^*$-fixed points $p_0, p_1$, or a covering of $\mathbb{P}^1$ fully ramified over $p_0$ and $p_1$. Associate a labeled graph $\Gamma$ to the $\mathbb{C}^*$-fixed point $\hat{f} : (C; x_1, \ldots, x_{\ell(\mu)}, z_1, \ldots, z_n) \to \mathbb{P}^1[m]$ as follows:

1. For each connected component $C_v$ of $\hat{f}^{-1}(\{p_0, p_1\})$, assign a vertex $v$, a label $g(v)$ which is the arithmetic genus of $C_v$, and a label $i(v) = \begin{cases} 0, & \text{if } \hat{f}(C_v) = p_0 \\ 1, & \text{if } \hat{f}(C_v) = p_1 \end{cases}.$

Denote by $V(\Gamma)^{(k)}$ the set of vertices $v$ with $i(v) = k$, for $k = 0, 1$. The set $V(\Gamma)$ of vertices of the graph $\Gamma$ will then be the disjoint union of $V(\Gamma)^{(0)}$ and $V(\Gamma)^{(1)}$.

2. Assign an edge $e$ to each rational irreducible component $C_v$ of $C$ such that $\hat{f}|_{C_v}$ is not a constant map. Then $\hat{f}|_{C_v}$ is fully ramified over $p_0$ and $p_1$ with degree $d(e)$. Let $E(\Gamma)$ denote the set of edges of $\Gamma$.

3. The set of flags is given by $F(\Gamma) = \{(v, e) : v \in V(\Gamma), e \in E(\Gamma), C_v \cap C_e \neq \emptyset\}$.

4. For each $v \in V(\Gamma)$, define $d(v) = \sum_{(v, e) \in F(\Gamma)} d(e)$ and let $\nu(v)$ be the partition of $d(v)$ determined by $\{d(e) : (v, e) \in F(\Gamma)\}$. In case $m > 0$, we assign an additional label for each $v \in V(\Gamma)^{(1)}$: let $\mu(v)$ be the partition of $d(v)$ determined by the ramification of $\hat{f}|_{C_v} : C_v \to \mathbb{P}^1(m)$ over $p_1(m)$.

Let $G_\chi(\mathbb{P}^1, \mu)$ be the set of all the graphs associated to the $\mathbb{C}^*$-fixed points in $\overline{\mathcal{M}}_{g,n}^\chi(\mathbb{P}^1, \mu)$. We now describe the set of fixed points associated to a given graph $\Gamma \in G_\chi(\mathbb{P}^1, \mu)$.

- Case $m = 0$: Any $\mathbb{C}^*$-fixed point in $\overline{\mathcal{M}}_{g,n}^\chi(\mathbb{P}^1, \mu)$ which is represented by a morphism to $\mathbb{P}^1[0] = \mathbb{P}^1$ is associated to a graph $\Gamma_0 \in G_\chi^0(\mathbb{P}^1, \mu)$ such that

$$ V(\Gamma_0)^{(0)} = \{v_1^0, \ldots, v_k^0\}, \quad g(v_i^0) = g_i, \quad k \in \mathbb{N}, \quad \sum_i (2 - 2g_i) = \chi $$

$$ V(\Gamma_0)^{(1)} = \{v_1^\infty, \ldots, v_{\ell(\mu)}^\infty\}, \quad g(v_1^\infty) = \cdots = g(v_{\ell(\mu)}^\infty) = 0, $$

$$ E(\Gamma_0) = \{e_1, \ldots, e_{\ell(\mu)}\}, \quad d(e_i) = \mu_i \text{ for } i = 1, \ldots, \ell(\mu) $$

$$ j(v) = \text{number of } z_i \text{'s mapped to } v, \quad \sum_{V(\Gamma_0)^{(0)}} j(v) = n $$

The two end-points of the edge $e_i$ are $v_i^0$ and $v_i^\infty$ for some $1 \leq j \leq k$. Let $\mu(v_i^0) = \{\mu_j|e_j \text{ has } v_i^0 \text{ as an endpoint}\}$. Define

$$ \overline{\mathcal{M}}_{\Gamma_0} = \prod_{1 \leq i \leq k} \overline{\mathcal{M}}_{g_i, \mu(v_i^0) + j(v_i^0)} $$

where we take $\overline{\mathcal{M}}_{0,1} = \overline{\mathcal{M}}_{0,2} = \overline{\mathcal{M}}_{1,0} = \{\text{pt}\}$, then there is a morphism $i_{\Gamma_0} : \overline{\mathcal{M}}_{\Gamma_0} \to \overline{\mathcal{M}}_{g,n}^\chi(\mathbb{P}^1, \mu)$ whose image is the fixed locus $F_{\Gamma_0}$ associated to $\Gamma_0$. Hence $i_{\Gamma_0}$ induces an isomorphism $\overline{\mathcal{M}}_{\Gamma_0}/A_{\Gamma_0} \cong F_{\Gamma_0}$ where $A_{\Gamma_0}$ is the
automorphism group of any morphism associated to the graph $\Gamma_0$, which can be obtained from the short exact sequence

$$1 \rightarrow \prod_{i=1}^{l(\mu)} \mathbb{Z}_{\mu_i} \rightarrow A_{\Gamma_0} \rightarrow \text{Aut}(\Gamma_0) \rightarrow 1$$

The virtual dimension of $F_{\Gamma_0}$ with only stable vertices is

$$d_{\Gamma_0} = \frac{3}{2} \chi + l(\mu) + n$$

For any vertex $v \in V(\Gamma_0)^{(0)}$, introduce multiplicity of $v$, $m(v)$, as follows:

$$m(v) = |\{ w \in V(\Gamma_0)^{(0)} \mid g(v) = g(w), \mu(v) = \mu(w), j(v) = j(w)\}|$$

Pick one representatives from each group of identical vertices $\{v_1, \ldots, v_l\}$ so that $\sum m(v_k) = |V(\Gamma_0)^{(0)}|$. Denote by $m(v_k) = m_k$ and now we can find the order of automorphism group of any $\Gamma_0 \in G_{\chi,n}(\mathbb{P}^1, \mu)$ to be:

$$|\text{Aut} \Gamma_0| = \prod_k m_k! \cdot j(v_k)! \cdot \text{Aut}(\mu(v_k))^{m_k}$$

- **Case $m > 0$:** For any given $f : (C; x_1, \ldots, x_{l(\mu)}, z_1, \ldots, z_n) \rightarrow \mathbb{P}^1[m]$, consider $f_0$ and $f_\infty$ which are defined as follows:
  - $f_0$: Let $C_0 = f^{-1}(p_0^{(0)})$, $\{y_1, \ldots, y_{l(\nu)}\} = f^{-1}(p_1^{(0)})$, and $\{z_1, \ldots, z_{n_0}\}$ mapped to $C_0$. Then $f_0 : (C_0; y_1, \ldots, y_{l(\nu)}, z_1, \ldots, z_{n_0}) \rightarrow \mathbb{P}^1$ corresponds to the case of $m = 0$.
  - $f_\infty$: Let $C_\infty = f^{-1}(\mathbb{P}^1(m))$ and $f_\infty = f |_{C_\infty}$, then $f_\infty$ corresponds to an element of $\mathcal{M}_{\Gamma}^{(1)}$ corresponding to the graph $\Gamma$ described below.

Classify the vertices of the graph $\Gamma$ as follows:

$$V^I(\Gamma)^{(0)} = \{ v \in V(\Gamma)^{(0)} : r_0(v) = -1 \},$$

$$V^{II}(\Gamma)^{(i)} = \{ v \in V(\Gamma)^{(i)} : r_i(v) = 0 \}, \quad \text{for } i = 0, 1$$

$$V^S(\Gamma)^{(i)} = \{ v \in V(\Gamma)^{(i)} : r_i(v) > 0 \}, \quad \text{for } i = 0, 1$$

where $r_0(v) = 2g(v) - 2 + \text{val}(v) + j(v)$, for $v \in V(\Gamma)^{(0)}$,

$$r_1(v) = 2g(v) - 2 + l(\mu(v)) + l(\nu(v)) + j(v), \quad \text{for } v \in V(\Gamma)^{(1)}$$

Define $\mathcal{M}_{\Gamma} = \mathcal{M}_{\Gamma}^{(0)} \times \mathcal{M}_{\Gamma}^{(1)}$, where $\mathcal{M}_{\Gamma}^{(0)} = \prod_{v \in V^S(\Gamma)^{(0)}} \mathcal{M}_{g(v), \text{val}(v)+j(v)}$ and $\mathcal{M}_{\Gamma}^{(1)}$ is the moduli space of morphisms $\hat{f} : \hat{C} \rightarrow \mathbb{P}^1(m)$ such that

1. $\hat{C} = \bigsqcup_{v \in V(\Gamma)^{(1)}} C_v$
2. For each $v \in V(\Gamma)^{(1)}$, $(C_v, x_{v,1}, \ldots, x_{v,l(\mu(v))}, y_{v,1}, \ldots, y_{v,l(\nu(v))}, z_1, \ldots, z_{j(v)})$ is a prestable curve of genus $g(v)$ with $l(\mu(v)) + l(\nu(v)) + j(v)$ ordered marked points.
(3) As Cartier divisors,
\[(\hat{f} \mid_{C_v})^{-1}(p_1^{(0)}) = \sum_{j=1}^{l(\nu)} \nu(v)_j y_{v,j}, \quad (\hat{f} \mid_{C_v})^{-1}(p_1^{(m)}) = \sum_{i=1}^{l(\mu)} \mu(v)_i x_{v,i}\]
The morphism \((\hat{f} \mid_{C_v})^{-1}(E) \to E\) is of degree \(d(v)\) for each irreducible component \(E\) of \(\mathbb{P}^1(m)\).

(4) The automorphism group of \(\hat{f}\) is finite. Here, an automorphism of \(\hat{f}\) consists of an automorphism of the domain curve \(\hat{C}\) and an automorphism of the pointed curve \((\mathbb{P}^1(m), p_1^{(0)}, p_1^{(m)})\), which is an element of \((\mathbb{C}^*)^m\).

There is a morphism \(i_\Gamma : \overline{\mathcal{M}}_\Gamma \to \overline{\mathcal{M}}_{\chi, \mu}(\mathbb{P}^1, \mu)\) whose image is the fixed locus \(F_\Gamma\) associated to the graph \(\Gamma\). Hence \(i_\Gamma\) induces an isomorphism \(\overline{\mathcal{M}}_\Gamma / A_\Gamma \cong F_\Gamma\) where \(A_\Gamma\) is the automorphism group of any morphism associated to the graph \(\Gamma\), which can be obtained from the short exact sequence
\[1 \to \prod_{e \in E(\Gamma)} \mathbb{Z}_{d(e)} \to A_\Gamma \to \text{Aut}(\Gamma) \to 1\]

The virtual dimension of \(\overline{\mathcal{M}}_\Gamma^{(1)}\) is given by \(d_\Gamma^{(1)} = \left(\sum_{v \in V(\Gamma^{(1)})} r_1(v)\right) - 1\). Use the identities
\[-\chi = -\chi_0 - \chi_\infty + 2l(\nu), \quad g = \sum_{v \in V(\Gamma)} g(v) - |V(\Gamma)| + |E(\Gamma)| + 1\]
to get the virtual dimension of \(F_\Gamma\):
\[d_\Gamma = d_\Gamma^{(0)} + d_\Gamma^{(1)} = -\frac{3}{2} \chi_0 + l(\nu) - \chi_\infty + l(\mu) + l(\nu) - 1 + n_\infty\]
\[= \sum_{v \in V_S(\Gamma)^{(0)}} (3g(v) - 3 + \text{val}(v)) + \left(\sum_{v \in V(\Gamma)^{(1)}} r_1(v)\right) - 1 + n_\infty\]
\[= 2g - 3 + l(\mu) + \sum_{v \in V_S(\Gamma)^{(0)}} (g(v) - 1) + |V^f(\Gamma)^{(0)}| + n_\infty\]
\[= -\frac{1}{2} \chi_0 - \chi + l(\mu) - 1 + n_\infty\]

By similar observation as in the case of \(m = 0\), we find the order of automorphism group of any given \(\Gamma \in G_\chi^\infty(\mathbb{P}^1, \mu)\) to be:
\[(8) \quad |\text{Aut } \Gamma| = |\text{Aut } \Gamma_0| |\text{Aut } \mu| \left(\prod_{n \in V(\Gamma)^{(1)}} n_k!\right) \left(\prod_{j(v) \in V(\Gamma)^{(1)}} j(v)! |\text{Aut } \mu(v)| |\text{Aut } \nu(v)|\right)\]

In this case, \(n_k\) is the multiplicity of vertices in \(V(\Gamma)^{(1)}\) with identical \(\mu(v), \nu(v), j(v),\) and \(g(v)\). The presence of automorphism group of \(\mu\) is due to the fact that we can exchange two marked points with same ramification type \(\mu_i\) without changing the type of corresponding graph.
4. Computation of $e_T(N_T^{\text{vir}})$

In this section, I will summarize the computations of $e_T(N_T^{\text{vir}})$ in [12] which will be needed for localization computation. Denote by $(\omega)$ the 1-dimensional representation of $\mathbb{C}^*$ given by $\lambda \cdot z = \lambda^w z$ for $\lambda \in \mathbb{C}^*$, $z \in \mathbb{C}$. For a given graph $\Gamma \in G_{g,n}(\mathbb{P}^1, \mu)$, let

\begin{equation}
[f : (C, x_1, \cdots, x_l(\mu), z_1, \cdots, z_n) \rightarrow \mathbb{P}^1[m]]
\end{equation}

be a fixed point of the $\mathbb{C}^*$-action on $\overline{M}_{\chi,n}(\mathbb{P}^1, \mu)$ associated to $\Gamma$. Given a flag $(v, e) \in F(\Gamma)$, denote by $q_{(v,e)} \in C$ the node at which $C_v$ and $C_e$ intersect. Also let $\psi_{(v,e)}$ denote the first Chern class of the cotangent line bundles over $\overline{M}_\Gamma$, i.e. the fiber at the fixed point $\Box$ is given by $T^*_{q_{(v,e)}}C_v$. The Euler class $e_T(N_T^{\text{vir}})$ is given by:

\[
\frac{1}{e_T(N_T^{\text{vir}})} = \frac{e_T(\hat{T}^2)}{e_T(T^1)}
\]

where $\hat{V}$ denotes the moving part of any vector bundle $V$, and $T^1, T^2$ are the tangent space and the obstruction space of $\overline{M}_{\chi,n}(\mathbb{P}^1, \mu)$, respectively. Here, $T^1$ and $T^2$ can be computed through the following two exact sequences [12]:

\[
0 \rightarrow \text{Ext}^0(\Omega_C(D), \mathcal{O}_C) \rightarrow H^0(\mathcal{D}^*) \rightarrow T^1 \rightarrow \text{Ext}^1(\Omega_C(D), \mathcal{O}_C) \rightarrow H^1(\mathcal{D}^*) \rightarrow T^2 \rightarrow 0
\]

where $0 \rightarrow H^0(\mathcal{D}^*) \rightarrow H^0(C, f^*(\omega_{\mathbb{P}^1[m]}(\log p_1^{(m)})))^\vee \rightarrow H^0(\mathcal{D}^*) \rightarrow \bigoplus_{i=0}^{m-1} H^0_{\text{et}}(R^*_i)

\rightarrow H^1(\mathcal{D}^*) \rightarrow \bigoplus_{i=0}^{m-1} H^1_{\text{et}}(R^*_i) \rightarrow 0

where $\omega_{\mathbb{P}^1[m]}$ is the dualizing sheaf of $\mathbb{P}^1[m]$, $D = x_1 + \cdots + x_l(\mu)$ is the branch divisor, and for $n_l$ is the number of nodes over $p_1^{(l)}$:

\[
H^0_{\text{et}}(R^*_i) \cong \bigoplus_{q \in f^{-1}(p_i^{(l)})} T_q(f^{-1}(\mathbb{P}^1_I)) \cong \mathbb{C}^{\otimes n_l}
\]

\[
H^1_{\text{et}}(R^*_i) \cong (T_{p_i^{(l)}}\mathbb{P}^1_I \otimes T_{p_i^{(l)}}\mathbb{P}^1_{I+1})^{\otimes(n_l-1)}
\]

Recall the map $\pi[m] : \mathbb{P}^1[m] \rightarrow \mathbb{P}^1$, and observe that for $\hat{f} = \pi[m] \circ f$ we have

\[
f^*(\omega_{\mathbb{P}^1[m]}(\log p_1^{(m)}))^\vee \cong \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)
\]

Let $F_\Gamma$ be the set of fixed points associated to $\Gamma \in G_{\chi,n}(\mathbb{P}^1, \mu)$ and assume that

\[
[f : (C, x_1, \cdots, x_l(\mu), z_1, \cdots, z_n) \rightarrow \mathbb{P}^1[m]] \in F_\Gamma \subset \overline{M}_{\chi,n}(\mathbb{P}^1, \mu)
\]
The $\mathbb{C}^*$-action on $\overline{\mathcal{M}}_{g,n}(\mathbb{P}^1, \mu)$ induces $\mathbb{C}^*$-actions on

$$\text{Ext}^0(\Omega_C(D), \mathcal{O}_C), \quad H^0(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)),$$

$$\text{Ext}^1(\Omega_C(D), \mathcal{O}_C), \quad H^1(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)),$$

$$\bigoplus_{l=0}^{m-1} H^0_{\text{et}}(\mathbb{R}^*_l),$$

$$\bigoplus_{l=0}^{m-1} H^1_{\text{et}}(\mathbb{R}^*_l).$$

The moving part of each of these groups form vector bundles over $\overline{\mathcal{M}}_g$. We will use the same notation $\hat{\cdot}$ to denote the induced vector bundles. In particular,

$$\bigoplus_{l=0}^{m-1} H^0_{\text{et}}(\mathbb{R}^*_l) = 0, \quad \text{and}$$

$$\bigoplus_{l=0}^{m-1} H^1_{\text{et}}(\mathbb{R}^*_l) = \begin{cases} 0, & \text{if } m = 0 \\ H^1_{\text{et}}(\mathbb{R}^*_0) = \left( T_{p_1(0)} \mathbb{P}^1 \otimes T_{p_2(0)} \mathbb{P}^1 \right)^{\oplus(n_0 - 1)}, & \text{if } m > 0 \end{cases}$$

Hence we have;

$$\frac{1}{e_T(N^\text{vir}_g)} = e_T(T^2) - e_T(T^1 - e_T(\text{Ext}^0(\Omega_C(D), \mathcal{O}_C))e_T(H^1(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)))e_T(\bigoplus_{l=0}^{m-1} H^1_{\text{et}}(\mathbb{R}^*_l)))$$

$$e_T(H^0(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)))e_T(\text{Ext}^1(\Omega_C(D), \mathcal{O}_C)).$$

**Case $m = 0$:**

For each $v \in V(\Gamma_0)$ and $\mu_{v,1}, \cdots, \mu_{v,l(l(v))}$ the ramification type in the vertex $v$, we have under the convention to write $\mu_{v,2} = \infty$ when $g(v) = 0, l(l(v)) = l(e(v)) = 1$;

$$\bigoplus_{l=0}^{m-1} H^1_{\text{et}}(\mathbb{R}^*_l)_v = 0$$

$$\text{Ext}^0(\Omega_C(D), \mathcal{O}_C)_v = \begin{cases} \left( \frac{1}{\mu_{v,1}} \right), & \text{if } v \in I \\ 0, & \text{if } v \in II \text{ or } S \end{cases}$$

$$\text{Ext}^1(\Omega_C(D), \mathcal{O}_C)_v = \begin{cases} 0, & \text{if } v \in I \\ \left( \frac{1}{\mu_{v,1}} + \frac{1}{\mu_{v,2}} \right), & \text{if } v \in II \\ \bigoplus_{i=1}^{l(l(v))} T_{q(v,e,v,i)} C_v \otimes T_{q(v,e,v,i)} C_v, & \text{if } v \in S \end{cases}.$$
Hence we can compute their contributions to be:

\[ e_T(\bigoplus_{i=0}^{m-1} H^1_{c_1}(R^*_i)_{v}) = 1 \]

\[ e_T(\text{Ext}^0(\Omega_C(D), O_C)_v) = \begin{cases} \frac{u}{\mu_{v,1}}, & \text{if } v \in I \\ 1, & \text{if } v \in \text{II or S} \end{cases} \]

\[ e_T(\text{Ext}^1(\Omega_C(D), O_C)_v) = \begin{cases} 1, & \text{if } v \in I \\ \frac{u}{\mu_{v,1}} + \frac{u}{\mu_{v,2}}, & \text{if } v \in \text{II} \\ \prod_{i=1}^{l(\mu(v))} \left( \frac{u}{\mu_{v,i}} - \psi_{v,i} \right), & \text{if } v \in S \end{cases} \]

For the contributions from the rest, consider the normalization sequence when \( v \in S; \)

\[ 0 \to \hat{f}^*O_{\mathbb{P}^1}(1) \to (\hat{f} |_{C_v})^*O_{\mathbb{P}^1}(1) \oplus \bigoplus_{i=1}^{l(\mu(v))} (\hat{f} |_{\mu_{v,i}})^*O_{\mathbb{P}^1}(1) \to \bigoplus_{i=1}^{l(\mu(v))} O_{\mathbb{P}^1}(1)_{p_0} \to 0 \]

The corresponding long exact sequence becomes

\[ 0 \to H^0(C, \hat{f}^*O_{\mathbb{P}^1}(1)) \to H^0(C_v, (\hat{f} |_{C_v})^*O_{\mathbb{P}^1}(1)) \oplus \bigoplus_{i=1}^{l(\mu(v))} H^0(C_{v,i}, (\hat{f} |_{C_{v,i}})^*O_{\mathbb{P}^1}(1)) \to H^1(C, \hat{f}^*O_{\mathbb{P}^1}(1)) \]

\[ \to H^1(C_v, (\hat{f} |_{C_v})^*O_{\mathbb{P}^1}(1)) \oplus \bigoplus_{i=1}^{l(\mu(v))} H^1(C_{v,i}, (\hat{f} |_{C_{v,i}})^*O_{\mathbb{P}^1}(1)) \to 0 \]

and the representations of \( C^* \) are given by

\[ 0 \to H^0(C, \hat{f}^*O_{\mathbb{P}^1}(1)) \to H^0(C_v, O_{C_v}) \otimes (1) \oplus \bigoplus_{i=1}^{l(\mu(v))} \left( \bigoplus_{a=1}^{\mu_{v,i}} (\frac{a}{\mu_{v,i}}) \right) \to \]

\[ \bigoplus_{i=1}^{l(\mu(v))} (1) \to H^1(C, \hat{f}^*O_{\mathbb{P}^1}(1)) \to H^1(C_v, O_{C_v}) \otimes (1) \to 0 \]

Hence their ratio can computed as;

\[ \frac{e_T(H^1(C, \hat{f}^*O_{\mathbb{P}^1}(1)))}{e_T(H^0(C, \hat{f}^*O_{\mathbb{P}^1}(1)))} = \prod_v \left[ \Lambda^v_{g(v)}(u)u^{l(\mu(v))-1} \prod_{i=1}^{l(\mu(v))} \left( \frac{\mu_{v,i}}{\mu_{v,i}} - u^{-\mu_{v,i}} \right) \right] \]

which also works for the case of \( v \in \text{I} \) or \( v \in \text{II}. \)

**Case** \( m > 0 \): Let \( \psi^i \) be the first Chern class of the line bundle over \( \overline{\mathcal{M}}_1^{(1)} \) whose fiber at \( \left[ f : \hat{C} \to \mathbb{P}^1(m) \right] \) is \( T^*_{\hat{p}^{-1}}(m) \). So \( \psi^i = \nu_{v,i} \psi_{(v,e,v,i)} \) for \( v \in V(\Gamma)^{(1)}, (v,e,v,i) \in F. \)
By similar observation as in the case of \( m = 0 \), we can find that

\[
\bigoplus_{l=0}^{m-1} H^l_{et}(R^*_l) = \left( T_{p_1^{(0)}} \mathbb{P}^1_{(0)} \otimes T_{p_1^{(0)}} \mathbb{P}^1_{(1)} \right)^{|E(G)|-1}
\]

\[
\text{Ext}^0(\Omega_C(D), O_C)_v = \begin{cases} \\
(\frac{1}{\nu_{v,1}}), & \text{if } v \in I \\
0, & \text{if } v \in II \text{ or } S \\
\end{cases}
\]

\[
\text{Ext}^1(\Omega_C(D), O_C)_v = \begin{cases} \\
0, & \text{if } v \in I \\
(\frac{1}{\nu_{v,1}} + \frac{1}{\nu_{v,2}}), & \text{if } v \in II \\
\bigoplus_{i=1}^{\nu(v)} T_{q(v,e_{v,i})} C_v \otimes T_{q(v,e_{v,i})} C_{v,i}, & \text{if } v \in S \text{ or } T \\
\end{cases}
\]

Hence we can compute their contributions to be;

\[
e_T\left( \bigoplus_{l=0}^{m-1} H^l_{et}(R^*_l) \right) = \left( -u - \psi^f \right)^{|E(G)|-1}
\]

\[
e_T(\text{Ext}^0(\Omega_C(D), O_C)_v) = \begin{cases} \\
u_{v,1}, & \text{if } v \in I \\
1, & \text{if } v \in II \text{ or } S \\
\end{cases}
\]

\[
e_T(\text{Ext}^1(\Omega_C(D), O_C)_v) = \begin{cases} \\
1, & \text{if } v \in I \\
\frac{1}{\nu_{v,1}} + \frac{u}{\nu_{v,2}}, & \text{if } v \in II \\
\prod_{i=1}^{\nu(v)} \left( \frac{u}{\nu_{v,i}} - \psi_{v,i} \right), & \text{if } v \in S \\
\prod_{i=1}^{\nu(v)} \left( \frac{u}{\nu_{v,i}} - \psi_{v,i} \right), & \text{if } v \in T \\
\end{cases}
\]

Also consider the following normalization sequence

\[
0 \to \hat{f}^* O_{\mathbb{P}^1}(1) \to \bigoplus_{S \cup T} \left( \hat{f} \mid c_v \right)^* O_{\mathbb{P}^1}(1) \oplus \bigoplus_{e \in E(G)} \left( \hat{f} \mid c_e \right)^* O_{\mathbb{P}^1}(1) \to \bigoplus_{I_H} O_{\mathbb{P}^1}(1)_{p_0} \oplus \bigoplus_{S} \left( \bigoplus_{(v,e) \in F} O_{\mathbb{P}^1}(1)_{p_0} \right) \oplus \bigoplus_{T} \left( \bigoplus_{(v,e) \in F} O_{\mathbb{P}^1}(1)_{p_1} \right) \to 0
\]

and the corresponding long exact sequence

\[
0 \to H^0(C, \hat{f}^* O_{\mathbb{P}^1}(1)) \to \bigoplus_{S \cup T} H^0(C_v, \left( \hat{f} \mid c_v \right)^* O_{\mathbb{P}^1}(1)) \oplus \bigoplus_{e \in E(G)} H^0(C_e, \left( \hat{f} \mid c_e \right)^* O_{\mathbb{P}^1}(1)) \to \bigoplus_{I_H} O_{\mathbb{P}^1}(1)_{p_0} \oplus \bigoplus_{S} \left( \bigoplus_{(v,e) \in F} O_{\mathbb{P}^1}(1)_{p_0} \right) \oplus \bigoplus_{T} \left( \bigoplus_{(v,e) \in F} O_{\mathbb{P}^1}(1)_{p_1} \right) \to H^1(C, \hat{f}^* O_{\mathbb{P}^1}(1)) \to \bigoplus_{S \cup T} H^1(C_v, \left( \hat{f} \mid c_v \right)^* O_{\mathbb{P}^1}(1)) \oplus \bigoplus_{e \in E(G)} H^1(C_e, \left( \hat{f} \mid c_e \right)^* O_{\mathbb{P}^1}(1)) \to 0
\]
The representations of $\mathbb{C}^*$ are given by

$$0 \rightarrow H^0(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)) \rightarrow \bigoplus_S H^0(C_v, \mathcal{O}_{C_v}) \otimes (1) \oplus \bigoplus_T H^0(C_v, \mathcal{O}_{C_v}) \otimes (0) \oplus \bigoplus_{e \in E(T)} \left( \bigoplus_{a=1}^{d(e)} \left( \frac{a}{d(e)} \right) \right).$$

from which we can compute their ratio to be;

$$\frac{e_T(H^1(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)))}{e_T(H^0(C, \hat{f}^*\mathcal{O}_{\mathbb{P}^1}(1)))} = \prod_{V((v))} \left[ \Lambda_{g(v)}(u) u^{l(v)-1} \right] \prod_{i=1}^{l(v)} \left[ \frac{u^{\nu_i}}{\nu_i!} u^{-\nu_i} \right].$$

After combining all the contributions, we find the following Feynman rules;

$$\frac{1}{e_T(N^{vir}_T)} = \left[ \prod_{i=1}^{l(\mu)} \frac{\mu_i^{\mu_i}}{\mu_i!} u^{-\mu_i} \right] \times \left[ \prod_i \frac{u}{\mu_{v,1}} \right] \times \left[ \prod_{I \in T} \frac{u}{\mu_{v,2}} \right] \times \left[ \prod_s \frac{\Lambda_{g(v)}(u)}{u} \left( \prod_{i=1}^{l(v)} \frac{u^{\nu_i}}{\nu_i!} u^{-\nu_i} \right) \right].$$

$$\frac{1}{e_T(N^{vir}_T)} = \frac{-\prod \nu_i}{u + \psi} \times \left[ \prod_{i=1}^{l(v)} \frac{\nu_i^{\nu_i}}{\nu_i!} u^{-\nu_i} \right] \times \left[ \prod_i \frac{u}{\nu_{v,1}} \right] \times \left[ \prod_{I \in T} \frac{u}{\nu_{v,2}} \right] \times \left[ \prod_s \frac{\Lambda_{g(v)}(u)}{u} \left( \prod_{i=1}^{l(v)} \frac{u^{\nu_i}}{\nu_i!} u^{-\nu_i} \right) \right].$$

5. Double Hurwitz Numbers

5.1. General Result of Hurwitz Numbers. Let $X$ be a Riemann surface of genus $h$. Given $n$ partitions $\eta^1, \ldots, \eta^n$ of $d$, denote by $H^X_d(\eta^1, \ldots, \eta^n)$ and $H^X_d(\eta^1, \ldots, \eta^n)$ the weighted counts of possibly disconnected and connected Hurwitz covers of type $(\eta^1, \ldots, \eta^n)$, respectively. The following Burnside formula is well known:

$$H^X_d(\eta^1, \ldots, \eta^n) = \sum_{|\rho| = d} \left( \frac{\dim R_{\rho}}{d!} \right)^{2-2h} \prod_{i=1}^{n} \left| C_{\eta^i} \right| \frac{\chi_{\rho}(C_{\eta^i})}{\dim R_{\rho}}.$$
5.2. Double Hurwitz Numbers. Consider a cover \( C \rightarrow \mathbb{P}^1 \) of genus \( g \), ramification type \( \nu, \mu \) at two points \( p_0 \) and \( p_1 \) respectively, and ramification type \((2)\) at \( r \) other points. By Riemann-Hurwitz formula, we have \( r = 2g - 2 + l(\nu) + l(\mu) \). Let \( \eta^1 = \cdots = \eta^r = (2) \) and introduce notations

\[
H^0_\nu(\nu, \mu) = H^0_\nu(\nu, \mu, \eta^1, \cdots, \eta^r)^\circ, \quad H^\bullet_\chi(\nu, \mu) = H^\bullet_\chi(\nu, \mu, \eta^1, \cdots, \eta^r)^\circ
\]

By applying Burnside formula, we obtain

\[
H^\bullet_\chi(\nu, \mu) = \sum_{|\xi|=d} \left( \left| C(2) \right| \frac{\chi(\xi(C(2)))}{\dim R_\xi} \right)^r \frac{\chi(\xi(C_\nu))}{z_\nu} \frac{\chi(\xi(C_\mu))}{z_\mu} = \sum_{|\xi|=d} f_\xi(2)^r \frac{\chi(\xi(C_\nu))}{z_\nu} \frac{\chi(\xi(C_\mu))}{z_\mu}
\]

Define generating series of Double Hurwitz numbers as follows:

\[
\Phi^0_{\nu, \mu}(\lambda) = \sum_{g \geq 0} H^0_g(\nu, \mu) \frac{\lambda^{2g - 2 + l(\nu) + l(\mu)}}{(2g - 2 + l(\nu) + l(\mu))!} \quad \Phi^0(\lambda; p_0^0, p_\infty^\infty) = \sum_{\nu, \mu} \Phi^0_{\nu, \mu}(\lambda) p_\nu^0 p_\mu^\infty
\]

\[
\Phi^\bullet_{\nu, \mu}(\lambda) = \sum_{\chi} H^\bullet_\chi(\nu, \mu) \frac{\lambda^{-\chi + l(\nu) + l(\mu)}}{(-\chi + l(\nu) + l(\mu))!} \quad \Phi^\bullet(\lambda; p_0^0, p_\infty^\infty) = 1 + \sum_{\nu, \mu} \Phi^\bullet_{\nu, \mu}(\lambda) p_\nu^0 p_\mu^\infty
\]

We will need the following initial value formula of Double Hurwitz numbers;

**Lemma 5.1.**

\[
\Phi^\bullet_{\nu, \mu}(0) = \frac{1}{z_\nu} \delta_{\nu, \mu}
\]

**Proof.** This is a direct consequence of the orthogonal relations for characters of \( S_d \)

\[
\sum_{\xi} \frac{\chi(\xi(C_\nu))}{z_\nu} = \delta_{\nu, \mu}, \quad \text{and} \quad \sum_{|\xi|=d} \frac{\chi(\xi(C_\nu))}{z_\nu} = z_\nu \delta_{\nu, \mu}
\]

\( \Box \)

5.3. Relating Double Hurwitz Numbers with Hodge Integrals. We can extend the notion of \( \mathbb{P}^1[m] \) to have bubble components on both directions. Denote by

\[
\mathbb{P}^1[m_0, m_\infty] = \mathbb{P}^1(-m_0) \cup \cdots \cup \mathbb{P}^1(-1) \cup \mathbb{P}^1(0) \cup \mathbb{P}^1(1) \cup \cdots \cup \mathbb{P}^1(m_\infty)
\]

As before, we call \( \mathbb{P}^1(0) \) the root component, \( \mathbb{P}^1_0(m_0) = \mathbb{P}^1(-m_0) \cup \cdots \cup \mathbb{P}^1(-1) \) the bubble component at \( 0 \), and \( \mathbb{P}^1_\infty(m_\infty) = \mathbb{P}^1(1) \cup \cdots \cup \mathbb{P}^1(m_\infty) \) the bubble component at \( \infty \). Define \( \overline{\mathcal{M}}_{\chi,n}(\mu^0, \mu^\infty) \) as the moduli space of morphisms \( f : C \rightarrow \mathbb{P}^1[m_0, m_\infty] \) such that

1. \( (C; x_1, \cdots, x_i(m_\infty), y_1, \cdots, y_i(m^0)) \) is a possibly-disconnected prestable curve with \( l(\mu^0) \) unordered marked points.
2. \( \chi = \sum (2 - 2g_i) \) where \( g_i \) is the genus of each connected component of \( C \).
3. As Cartier divisors, \( f^{-1}(p_0^{(m_0)}) = \sum_{j=1}^{l(\mu^0)} p_0^j y_j \), \( f^{-1}(p_1^{(m_\infty)}) = \sum_{i=1}^{l(\mu^\infty)} p_1^\infty x_i \)
4. The automorphism group of \( f \) is finite. Here, an automorphism of \( f \) consists of an automorphism of the domain curve \( C \) and automorphisms of the pointed curves \( (\mathbb{P}^1_0(m_0), p_0^{(-m_0)}, p_1^{(-1)}) \) and \( (\mathbb{P}^1_\infty(m_\infty), p_0^{(1)}, p_1^{(m_\infty)}) \), which is an element of \( (\mathbb{C}^*)^{m_0} \) and \( (\mathbb{C}^*)^{m_\infty} \), respectively.
We can extend the standard action \( t \cdot [z, w] = [tz, wz] \) on \( \mathbb{P}^1 \) to \( \mathbb{P}^1[m_0, m_\infty] \) by trivial action on the bubble components at 0 and \( \infty \). Then this action induces \( \mathbb{C}^* \)-action on \( \overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \mu^\infty) \). Let \( \pi[m_0, m_\infty] : \mathbb{P}^1[m_0, m_\infty] \to \mathbb{P}^1 \) be the projection which contracts both bubble components and \( f_\rho = \pi[m_0, m_\infty] \circ f \). Denote by \( \nu \) the partition of \( d = |\mu^0| = |\mu^\infty| \) by the degrees of \( f_\rho \) on each rational irreducible components. For any morphism \( f : (C, x_i, y_j) \to \mathbb{P}^1[m_0, m_\infty] \) which represents a fixed point of \( \overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \mu^\infty) \) under this action, one of the following four cases must hold:

- \( m_0 = m_\infty = 0 \): We have \( f = f_\rho \), \( \mu^0 = \mu^\infty = \nu \)
- \( m_0 = 0, m_\infty > 0 \): Let \( \chi_\infty = \sum (2 - 2g_\infty^i) \) where \( g_\infty^i \) is the genus of each connected component of \( f^{-1}(\mathbb{P}^1(m_\infty)) \). In this case, we have \( \mu^0 = \nu, \chi_\infty = \chi, \chi_0 = 2l(\nu) \)
- \( m_0 > 0, m_\infty = 0 \): Let \( \chi_0 = \sum (2 - 2g_j^0) \) where \( g_j^0 \) is the genus of each connected component of \( f^{-1}(\mathbb{P}^1_0(m_0)) \). In this case, we have \( \mu^\infty = \nu, \chi_0 = \chi, \chi_\infty = 2l(\nu) \).
- \( m_0 > 0, m_\infty > 0 \): We have \( \chi = \chi_0 + \chi_\infty - 2l(\nu) \).

Hence, we can see that \( \chi_0, \chi_\infty, \nu \) determines each connected component of \( \overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \mu^\infty) \). Consider the branch morphism

\[
\text{Br} : \overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \mu^\infty) \to \text{Sym}^{-\chi + l(\mu^0) + l(\mu^\infty)} \mathbb{P}^1 \cong \mathbb{P}^{-\chi + l(\mu^0) + l(\mu^\infty)}
\]

The Double Hurwitz numbers for possibly disconnected covers of \( \mathbb{P}^1 \) can be defined by

\[
H^*_\chi(\mu^0, \mu^\infty) = \frac{1}{|\text{Aut}(\mu^0)||\text{Aut}(\mu^\infty)|} \int_{\overline{\mathcal{M}}_\chi(\mu^0, \mu^\infty)} \text{Br}^*(H^{-\chi + l(\mu^0) + l(\mu^\infty)})
\]

under the assumption \( -\chi + l(\mu^0) + l(\mu^\infty) > 0 \) where \( H \in H^2(\mathbb{P}^{-\chi + l(\mu^0) + l(\mu^\infty)}) \) is the hyperplane class. We want to compute this integration by virtual localization. The connected components of \( \overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \mu^\infty) \) can be described as follows:

- \( m_0 = 0, m_\infty > 0 \) : \( \mathcal{F}(\nu; 2l(\nu), \chi) \cong (\overline{\mathcal{M}}_\chi(\mathbb{P}^1, \nu, \mu^\infty) // \mathbb{C}^*) / \prod \mathbb{Z}_{\mu^\infty} \)
- \( m_0 > 0, m_\infty = 0 \) : \( \mathcal{F}(\nu; \chi, 2l(\nu)) \cong (\overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \nu) // \mathbb{C}^*) / \prod \mathbb{Z}_{\mu^0} \)
- \( m_0 > 0, m_\infty > 0 \) : \( \mathcal{F}(\nu; \chi_0, \chi_\infty) \cong ((\overline{\mathcal{M}}_{\chi_0}(\mathbb{P}^1, \mu^0, \nu) // \mathbb{C}^*) \times (\overline{\mathcal{M}}_{\chi_\infty}(\mathbb{P}^1, \nu, \mu^\infty) // \mathbb{C}^*) ) / (\text{Aut}(\nu) \prod_{i=1}^{l(\nu)} \mathbb{Z}_{\nu_i}) \)

Let \( \mathcal{N}^\text{vir}_{\nu, \chi_0, \chi_\infty} \) be the pull-back of the virtual normal bundle of \( \mathcal{F}(\nu, \chi_0, \chi_\infty) \) in \( \overline{\mathcal{M}}_\chi(\mathbb{P}^1, \mu^0, \mu^\infty) \). By computations similar to those \( e_T(\mathcal{N}^\text{vir}_{\nu; 2l(\nu), \chi}) \), we obtain

\[
e_T(\mathcal{N}^\text{vir}_{\nu, 2l(\nu), \chi}) = \frac{-a_\nu}{u + \psi^{0}}, \quad \frac{1}{e_T(\mathcal{N}^\text{vir}_{\nu, \chi, 2l(\nu)})} = \frac{a_\nu}{u - \psi^{\infty}}, \quad \frac{1}{e_T(\mathcal{N}^\text{vir}_{\nu, \chi_0, \chi_\infty})} = \frac{-A_\nu}{u + \psi^{0}} \times \frac{a_\nu}{u - \psi^{\infty}}
\]
where $\psi^0$ and $\psi^\infty$ are the first Chern classes of the cotangent line bundle $T^*_{\eta_0} \mathbb{P}^1[m_0, m_\infty]$ and $T^*_{\eta_1} \mathbb{P}^1[m_0, m_\infty]$, respectively. Let $r = -\chi + l(\mu^0) + l(\mu^\infty)$ and observe that

$$Br(F(\nu; \chi_0, \chi_\infty)) = (-\chi_\infty + l(\nu) + l(\mu^\infty))p_1 + (-\chi_0 + l(\mu^0) + l(\nu))p_0 \in \mathbb{P}^r$$

$$i^*_{\nu, \chi_0, \chi_\infty} Br^* \left( \prod_{k=1}^r (H - w_k) \right) = \left( \prod_{k=1}^r (-\chi_0 + l(\mu^0) + l(\nu) - w_k) \right) u^r$$

By taking special values for $w$:

$$w = (0, 1, \ldots, -\chi + l(\mu^0) + l(\mu^\infty) - 1) \quad \text{and} \quad w = (1, 2, \ldots, -\chi + l(\mu^0) + l(\mu^\infty))$$

, we can compute to obtain

$$H^\bullet(\mu^0, \mu^\infty) \over (-\chi + l(\mu^0) + l(\mu^\infty))! = \frac{1}{|\text{Aut}(\mu^0)| |\text{Aut}(\mu^\infty)|} \int_{[\overline{M}(\mathbb{P}^1, \mu^0, \mu^\infty)/\mathbb{C}^*]}^{\text{vir}} (\psi^0)^{-\chi + l(\mu^0) + l(\mu^\infty) - 1}$$

$$= \frac{1}{|\text{Aut}(\mu^0)| |\text{Aut}(\mu^\infty)|} \int_{[\overline{M}(\mathbb{P}^1, \mu^0, \mu^\infty)/\mathbb{C}^*]}^{\text{vir}} (\psi^\infty)^{-\chi + l(\mu^0) + l(\mu^\infty) - 1}$$

## 6. Recursion Formula

In this section, I will summarize the results from previous sections to prove the following recursion formula on Hodge integrals. Let $e = (k_1, \ldots, k_n)$ be a partition where $k_i$'s are allowed to be zero.

**Theorem 6.1.** For any partition $\mu$ and $e$ with $|e| < |\mu| + l(\mu) - \chi$, we have

$$\left[\lambda^{l(\mu) - \chi} \right] \sum_{|\nu| = |\mu|} \Phi_{\mu, \nu}^* (-\lambda) z_\nu D_{\nu, e}^* (\lambda) = 0$$

where the sum is taken over all partitions $\nu$ of the same size as $\mu$. 
Here \([\lambda^n]\) means taking the coefficient of \(\lambda^n\). Let me first introduce some notations;

\[
D_{g,\nu,\varepsilon} = \begin{cases} 
\frac{\nu_1^{\nu_1-2}}{\nu_1!}, & \text{if } (g, l(\nu) + l(e)) = (0, 1) \\
\frac{1}{|\text{Aut } \nu|} \frac{\nu_1^{\nu_1} \nu_2^{\nu_2}}{\nu_1! \nu_2!} \frac{1}{\nu_1 + \nu_2}, & \text{if } (g, l(\nu), l(e)) = (0, 2, 0) \\
\frac{\nu_1^{\nu_1}}{\nu_1!} \sum_{k=0}^{\varepsilon_1} \frac{1}{\nu_1^{1+k}} \binom{\varepsilon_1}{k}, & \text{if } (g, l(\nu), l(e)) = (0, 1, 1) \\
\frac{1}{l(e)! |\text{Aut } \nu|} \left[ \prod_{i=1}^{l(\nu)} \frac{\nu_i^{\nu_i}}{\nu_i!} \right] \int_{\mathcal{M}_{g, l(\nu) + l(e)}} \frac{\Lambda_g^\nu(1) \prod_{j=1}^{l(e)} (1 - \psi_j)^{e_j}}{\prod_{i=1}^{l(\nu)} (1 - \nu_i \psi_i)}, & \text{otherwise}
\end{cases}
\]

\[
D(\lambda, p, q) = \sum_{|\nu| \geq 0} \sum_{g \geq 0} \lambda^{2g-2+l(\nu)} p_\nu q_\nu D_{g, \nu}
\]

\[
D^*(\lambda, p, q) = \exp(D(\lambda, p, q)) = \sum_{|\nu| \geq 0} \lambda^{-\chi + l(\nu)} p_\nu q_\nu D^*_\chi p_e = \sum_{|\nu| \geq 0} p_\nu q_\nu D^*_\nu(\lambda)
\]

where \(p_i, q_j\)’s are formal variables with \(p_\nu = p_{\nu_1} \times \cdots \times p_{\nu_{l(\nu)}}\) and \(q_e = q_{\varepsilon_1} \times \cdots \times q_{\varepsilon_{l(e)}}\).

**Proof.** For any given \(\mu\) and \(\chi\) such that \(|\mu| + l(\mu) > \chi\), applying the localization formula to the class \(\prod_{j=1}^n \psi_j^{k_j} e^* \mu H\) where \(H\) is the hyperplane class of \(\mathbb{P}^1\) yields;

\[
0 = \int_{\mathcal{M}_{\chi,n}(\mathbb{P}^1, \mu)} \prod_{j=1}^n \psi_j^{k_j} e^* \mu H \quad \text{since } \deg \prod_{j=1}^n \psi_j^{k_j} e^* \mu H < \dim \mathcal{M}_{\chi,n}(\mathbb{P}^1, \mu)
\]

\[
= \sum_{\Gamma_0 \in \mathcal{G}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{1}{|A_{\Gamma_0}|} \int_{\mathcal{M}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{\prod_{j=1}^n (u - \psi_j)^{k_j} e^* \mu H_T}{e_T(\mathcal{N}_{\Gamma_0}^{\text{vir}})} + \sum_{\Gamma \in \mathcal{G}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{1}{|A_{\Gamma}|} \int_{\mathcal{M}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{\prod_{j=1}^n (u - \psi_j)^{k_j} e^* \mu H_T}{e_T(\mathcal{N}_{\Gamma}^{\text{vir}})}
\]

Here \(H_T\) is the lift of \(H\) to the equivariant hyperplane class. Choose \(H\) in such a way that \(H(0) = 0\) and \(H(\infty) = u\), then we have \(H_T(0) = u\) and \(H_T(\infty) = 0\). Also let \(u = 1\) for simplicity, then the formula reduces to:

\[
\sum_{\Gamma_0 \in \mathcal{G}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{1}{|A_{\Gamma_0}|} \int_{\mathcal{M}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{\prod_{j=1}^n (1 - \psi_j)^{k_j}}{e_T(\mathcal{N}_{\Gamma_0}^{\text{vir}})} + \sum_{\Gamma \in \mathcal{G}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{1}{|A_{\Gamma}|} \int_{\mathcal{M}_{\chi,n}(\mathbb{P}^1, \mu)} \frac{\prod_{j=1}^n (1 - \psi_j)^{k_j}}{e_T(\mathcal{N}_{\Gamma}^{\text{vir}})} = 0
\]

where \(\mathcal{G}_{\chi,n}^0\) and \(\mathcal{G}_{\chi,n}\) are the set of graphs corresponding to fixed locus with \(m = 0\) and \(m > 0\) with all the marked points \(z_1, \cdots, z_n\) concentrated on the vertices in
\[ V(\Gamma_0)^{(0)} \text{ and } V(\Gamma)^{(0)} \text{, respectively. We can compute the summand for } \Gamma_0 \text{ as follows:} \]
\[
\int_{\mathcal{M}_{\Gamma_0}} \frac{\prod_{j=1}^n (1 - \psi_j)^{k_j}}{e_T(\mathcal{N}_{\Gamma_0}^{\text{vir}})} = \left[ \prod_{i=1}^{\nu_i} \frac{\mu_i^{\nu_i}}{\nu_i!} \right] \times \left[ \prod_I \int_{\{pt\}} \frac{1}{\nu_{v,i}} \right] \times \left[ \prod_{II} \int_{\{pt\}} \frac{1}{\nu_{v,1} + 1/\nu_{v,2}} \right]
\times \left[ \prod_S \int_{\mathcal{N}_{\Gamma_0}^{\text{vir}}(g,v),l(\nu(v)) + l(e(v))} \frac{\Lambda^{\nu(v)} g(v)(1) \prod_{j=1}^{\nu(v)} (1 - \psi_{v,j})^{k_v,j}}{\prod_{i=1}^{l(v)(\nu)} (1 - \psi_{v,i})} \right]
\]
\[
= \left[ \prod_{I} \mu_{v,1}^{\nu_{v,1} - 1} \right] \times \left[ \prod_{II} \nu_{v,1} \nu_{v,2} \mu_{v,1}^{\nu_{v,2}} \mu_{v,2}^{\nu_{v,2}} \right] \times \left[ \prod_S \int_{\mathcal{N}_{\Gamma_0}^{\text{vir}}(g,v),l(\nu(v)) + l(e(v))} \frac{\Lambda^{\nu(v)} g(v)(1) \prod_{j=1}^{\nu(v)} (1 - \psi_{v,j})^{k_v,j}}{\prod_{i=1}^{l(v)(\nu)} (1 - \psi_{v,i})} \right]
\times \left[ \prod_{i=1}^{l(v)} \nu_{v,i} \right] \int_{\mathcal{N}_{\Gamma_0}^{(1)}} (\psi^t)^{-\chi_{\infty} + l(\mu) + l(\nu) - 1}
\]
\[
= \left[ \prod_{V(\Gamma_0)} z_{\nu(v),j}(v)! D_{g(v),\nu(v),e(v)} \right] \times \left[ \prod_{i=1}^{l(v)} \nu_{v,i} \right] \int_{\mathcal{N}_{\Gamma_0}^{(1)}} (\psi^t)^{-\chi_{\infty} + l(\mu) + l(\nu) - 1}
\]

And the integration over \( \mathcal{M}_{\Gamma_0}^{(1)} \) can be related to Double Hurwitz Numbers as follows:
\[
\text{Aut } \mu \text{ Aut } \nu \quad H_{\chi_{\infty}}^{\bullet}(\mu, \nu) = \left( -\chi_{\infty} + l(\mu) + l(\nu) \right)! \int_{\mathcal{N}_{\chi_{\infty}^{\bullet}(\mathbb{P}^1,\mu,\nu)/\mathbb{C}^*^{\text{vir}}}^{\text{vir}} (\psi^0)^{-\chi_{\infty} + l(\mu) + l(\nu) - 1}
\]
\[
= \frac{\left( -\chi_{\infty} + l(\mu) + l(\nu) \right)!}{\left[ \prod_{n_k!} \right] \left[ \prod_{V(\Gamma_0)} \text{Aut } \mu(\nu) \text{ Aut } \nu(v) \right] \int_{\mathcal{N}_{\Gamma_0}^{(1)}} (\psi^t)^{-\chi_{\infty} + l(\mu) + l(\nu) - 1}}
\]
since marked points in \( M^{(1)}_\Gamma \) are ordered. Recall that \(|A_{\Gamma_0}|\) and \(|A_{\Gamma}|\) are given by:

\[
|A_{\Gamma_0}| = \left( \prod_{i=1}^{l(\mu)} \mu_i \right) \prod_k m_k! \left( j(v_k)! \mid \text{Aut } \mu(v_k) \right)^{m_k}
\]

\[
|A_{\Gamma}| = \left( \prod_{i=1}^{l(\nu)} \nu_i \right) \left( \prod_k m_k! \left( j(v_k)! \mid \text{Aut } \nu(v_k) \right)^{m_k} \right) |\text{Aut } \mu| \times \left( \prod_{i=1}^{l(\nu)} \nu_i \right) \left( \prod_k m_k! \left( j(v_k)! \mid \text{Aut } \nu(v_k) \right)^{m_k} \right) |\text{Aut } \nu| \]

Also observe that

\[
\frac{\prod_{\Gamma \in \mathcal{G}_\chi,0} \sum_{V(\Gamma)} z_{\mu(\nu)} j(v)! \mathcal{D}_{g(\nu),\mu(v),e(v)}}{\prod_{i=1}^{l(\mu)} \mu_i \prod_k m_k! \left( j(v)! \mid \text{Aut } \mu(v_k) \right)^{m_k}} = \frac{1}{|V(\Gamma_0)|! (m_1, \ldots, m_t)} \prod_k \mathcal{D}_{g(v_k),\mu(v_k),e(v_k)}
\]

which is the coefficient of \( \lambda^{-\chi+l(\mu)} p_0 q_k \) in the expansion of \( \mathcal{D}\bullet(\lambda, p, q) \). Now the original equation can be simplified as follows;

\[
0 = \sum_{\Gamma_0 \in G_{\chi,0}(\mathbb{P}^1, \mu)} \frac{1}{|A_{\Gamma_0}|} \int_{X_{\Gamma_0}} \prod_{j=1}^{n} \left( 1 - \psi_j \right)^{k_j} e_T(N_{1\chi}^{\vir}) + \sum_{\Gamma \in G_{\chi,0}(\mathbb{P}^1, \mu)} \frac{1}{|A_{\Gamma}|} \int_{X_{\Gamma}} e_T(N_{1\chi}^{\vir})
\]

\[
= \sum_{\Gamma_0 \in G_{\chi,0}(\mathbb{P}^1, \mu)} \frac{1}{|V(\Gamma_0)|! (m_1, \ldots, m_t)} \prod_k \mathcal{D}_{g(v_k),\mu(v_k),e(v_k)}
\]

\[
+ \sum_{\Gamma \in G_{\chi,0}(\mathbb{P}^1, \mu)} \left[ \frac{1}{|V(\Gamma)(0)|! (m_1, \ldots, m_t)} \prod_k \mathcal{D}_{g(v_k),\nu(v_k),e(v_k)} \right.
\]

\[
\times (\lambda - \chi_\infty + l(\mu) + l(\nu) \prod_{i=1}^{l(\nu)} \nu_i) \left( \prod_{i=1}^{n} \mu_i \prod_{V(\Gamma)(1)} |\text{Aut } \mu| |\text{Aut } \nu| \right) \int_{X_{\Gamma}(1)} \psi_t (\lambda - \chi_\infty + l(\mu) + l(\nu) - 1)
\]

\[
= \sum_{\Gamma_0 \in G_{\chi,0}(\mathbb{P}^1, \mu)} \frac{1}{|V(\Gamma_0)|! (m_1, \ldots, m_t)} \prod_k \mathcal{D}_{g(v_k),\mu(v_k),e(v_k)}
\]

\[
+ \sum_{\Gamma \in G_{\chi,0}(\mathbb{P}^1, \mu)} \left[ \frac{1}{|V(\Gamma)(0)|! (m_1, \ldots, m_t)} \prod_k \mathcal{D}_{g(v_k),\nu(v_k),e(v_k)} \right.
\]

\[
\times (\lambda - \chi_\infty + l(\mu) + l(\nu) \sum_{\mu} z_{\mu(\nu)} \psi_t (\lambda - \chi_\infty + l(\mu) + l(\nu) - 1)
\]

\[
= \mathcal{D}^{\bullet}_{\chi,\mu,e} + \sum_{\nu} \left( -1 \right) ^{\chi_\infty + l(\mu) + l(\nu)} H_{\chi_\infty}(\mu, \nu) \mathcal{D}_{\chi_0,\nu,e}^{\bullet}
\]

\[
= \sum_{\nu} \left( -1 \right) ^{\chi_\infty + l(\mu) + l(\nu)} H_{\chi_\infty}(\mu, \nu) \mathcal{D}_{\chi_0,\nu,e}^{\bullet}
\]
where \( \chi = \chi_0 + \chi_\infty - 2l(\nu) \) and the initial value formula for Double Hurwitz numbers is used at the last equality. Summing over \( \chi \) yields that we have for all \( |e| + \chi < |\mu| + l(\mu) \)

\[
[\chi^{(\mu) - \chi} \sum_{|\nu| = |\mu|} \Phi_{\mu, \nu}(\lambda) z_\nu D_{\nu, e}(\lambda)] = 0
\]

\[\square\]

7. Computing Hodge Integrals with one \( \lambda \)-class

It is enough to consider the case of \( \mu = (d) \) for some positive integers \( d \) to compute all Hodge integrals with at-most one \( \lambda \)-class. And in this case, we have a closed formula for the Double Hurwitz Numbers as follows;

**Theorem 7.1.** ([2], Theorem 3.1.) Let \( r = r_{(d), \beta}^g \). For \( g \geq 0 \), and \( \beta \vdash d \) with \( n \) parts,

\[
H_{(d), \beta}^g = \lim_{t \to 0} \left[ t^{2g} \right] \prod_{k \geq 1} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} = \frac{r!}{2^g} \sum_{\lambda \vdash |\lambda|} \xi_{\lambda} S_{\lambda}
\]

where \( r = 2g - 1 + l(\beta) \) and \( c_1 = \text{number of 1's in } \beta \) \( - 1 \), \( c_k = \text{number of } k \text{'s in } \beta \) for \( k > 1 \). \( \left[ t^{2g} \right] \) means taking the coefficient of \( t^{2g} \).

Here the Double Hurwitz number is counted with multiplicity, hence in our notation it will read as follows

\[
H_{\chi_\infty}^*((d), \nu) = \frac{d-\chi_\infty + l(\nu)}{|\text{Aut } \nu|} \left[ t^{2g} \right] \prod_{k \geq 1} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k}
\]

And in this case, (10) can be written as

\[
\sum_{|e| = d, \chi_0, \chi_\infty} D_{\chi_0, \nu, e, z_\nu} (1 - \chi_\infty + 1 + l(\nu)) d^{\chi_\infty + l(\nu)} |\text{Aut } \nu| \left[ t^{2g} \right] \prod_{k \geq 1} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} = 0
\]

Now fix \( \nu \) and consider the case where there are \( m \) vertices in \( V(\Gamma)^{(0)} \). Then we have splitting of \( \chi_0, \nu, \) and \( e \) into \( \{g_1, \ldots, g_m\}, \{\nu(v_1), \ldots, \nu(v_m)\}, \) and \( \{e(v_1), \ldots, e(v_m)\} \) such that \( e(v_i) \)'s are allowed to be empty and

\[
\sum_{i=1}^{m} (2 - 2g_i) = \chi_0, \quad \bigcup_{i=1, \ldots, m} \nu(v_i) = \nu, \quad \bigcup_{i=1, \ldots, m} e(v_i) = e
\]

Each vertex will correspond to a certain Hodge integral on \( \overline{\mathcal{M}}_{g(\nu), l(\nu(e))} \) with dimension \( 3g(\nu) - 3 + l(\nu(e)) + l(e(\nu)) \). There are conditions on \( m, \chi, \chi_0, \chi_\infty, l(e(\nu)), \) and \( l(\nu(e)) \); let \( \overline{\mathcal{M}}(\nu) \) denote \( \sum_{w \neq e} g(w) \),

\[
m \leq l(\nu), \quad l(\nu(e)) \leq l(\nu) - m + 1, \quad \chi = \chi_0 + \chi_\infty - 2l(\nu), \quad l(e(\nu)) \leq l(e)
\]

\[
\chi_\infty \leq 2 \min \{ l((d)), l(\nu) \} = 2, \quad \chi_0 = \sum_{i=1}^{m} (2 - 2g(v_i)) = 2m - 2g(v) - 2\overline{\mathcal{M}}(\nu)
\]
From these conditions, we can deduce that

\[ 3g(v) = 3m - 3\bar{\gamma}(v) - \frac{3}{2}\chi_0 = 3m - 3\bar{\gamma}(v) - \frac{3}{2}\chi + \frac{3}{2}\chi_\infty - 3l(\nu) \]

\[ \leq \frac{3}{2}\chi + 3m - \frac{3}{2}(2l(\nu) - \chi_\infty) \]

where equality holds if and only if \( \bar{\gamma}(v) = \sum_{w \neq v} g(w) = 0 \). Now we can find the upper bound for the dimension of Hodge integral as follows.

\[ 3g(v) - 3 + l(\nu(v)) + l(e(v)) \leq -\frac{3}{2}\chi + 3m - 3 - \frac{3}{2}(2l(\nu) - \chi_\infty) + l(\nu(v)) + l(e) \]

\[ \leq \left[ 3g - 3 + l(e) + 1 \right] + \left[ 3m - 3 - 3l(\nu) + \frac{3}{2}\chi_\infty + l(\nu) - m \right] \]

\[ \leq \left[ 3g - 3 + l(e) + 1 \right] + \left[ 2(m - l(\nu)) + \frac{3}{2}(\chi_\infty - 2) \right] \]

\[ \leq 3g - 2 + l(e) \]

and the equality holds if and only if

\[ m = l(\nu), \quad \chi_\infty = 2, \quad e(w) = \emptyset \text{ for all } w \neq v, \quad g(w) = 0 \text{ for all } w \neq v \]

i.e. when each part of \( \nu \) is split into separate vertices, and all the marked points other than the ramification divisor as well as all genuses are concentrated on one vertex on the 0-th side. Now we can compute any Hodge integral with at-most one \( \lambda \)-class as follows: Say we want to compute Hodge integrals of the form

\[ \int_{\mathcal{M}_{g,n+1}} \psi_0^{k_0} \cdots \psi_n^{k_n} \lambda_j \]

where \( j + \sum_{i=0}^n k_i = 3g - 2 + n \). Assume \( 0 \leq k_0 \leq \cdots \leq k_n \) and let \( e = (k_1, \cdots, k_n) \) and \( \chi = 2 - 2g \). Then for any positive integer \( d \) such that \( d > \chi + |e| - 1 \), the recursion formula \( \text{(12)} \) expresses the top-dimensional Hodge integrals in terms of lower-dimensional Hodge integrals as follows:

\[ \sum_{|\nu| = d} \left[ \left( \prod_{i=1}^{l(\nu)} \nu_i^{(l(\nu) - 1)} \nu_i^{-1} \right) \frac{(-1)^{l(\nu) - 1} d^{l(\nu) - 2} g^{l(\nu)}}{n!} \sum_{i=1}^n \frac{\nu_i^2}{\text{Aut} \ \hat{\nu_i}} \int_{\mathcal{M}_{g,n+1}} \frac{\Lambda^\vee_g(1) \prod_{j=1}^n (1 - \psi_j)^{k_j}}{1 - \nu_i \psi_0} \right] \]

\[ = \text{terms consisting of lower dimensional Hodge integrals only...} \]
where $\text{Aut} \, \hat{\nu}_i$ is the automorphism group of the partition $\hat{\nu}_i = (\nu_1, \cdots, \nu_i, \cdots, \nu_{l(\nu)})$.
In this expression, the Hodge integral term expands to:
\[
\int_{\mathcal{M}_{g,n+1}} \frac{\Lambda_g^\nu (1)^n \prod_{j=1}^n (1 - \psi_j)^{k_j}}{1 - \nu_i \psi_0}
\]
\[
= \int_{\mathcal{M}_{g,n+1}} (1 - \lambda_1 + \lambda_2 + \cdots + (-1)^g \lambda_g) \left( \sum_{a_0=0}^\infty \nu_i^{a_0} \psi_0^{a_0} \right) \prod_{j=1}^n \left( \sum_{a_j=0} (-1)^{a_j} \left( \begin{array}{c} k_j \\ a_j \end{array} \right) \psi_j^{a_j} \right)
\]
\[
= \sum_{k + \sum a_j = 3g - 2 + n} (-1)^{3g-2+n} \left[ \nu_i^{a_0} \prod_{j=1}^n \left( \begin{array}{c} k_j \\ a_j \end{array} \right) \right] \int_{\mathcal{M}_{g,n+1}} \psi_0^{a_0} \times \cdots \times \psi_n^{a_n} \lambda_k
\]
Hence the previous expression can be written as
\[
(13) \sum_{k + \sum a_j = 3g - 2 + n} C_d(k, (a_j)) \int_{\mathcal{M}_{g,n+1}} \psi_0^{a_0} \times \cdots \times \psi_n^{a_n} \lambda_k = \text{lower dimensional terms...}
\]
where $C_d(k, (a_j))$ are constants defined as follows:
\[
C_d(k, (a_j)) = \sum_{|\nu|=d} \left[ \left( \prod_{i=1}^{l(\nu)} \nu_i^{\nu_i - 1} \right) \frac{(-1)^{l(\nu)-1} d^{l(\nu)-2}}{n!} \sum_{i=1}^{l(\nu)} \frac{\nu_i^2}{\text{Aut} \, \hat{\nu}_i} \left( \begin{array}{c} k_j \\ a_j \end{array} \right) \right]
\]
Now we have infinitely many linear relations of finitely many Hodge integrals of fixed dimension $3g - 2 + n$ with at-most one $\lambda$-class since the equation (13) holds for all positive integers $d$ such that $d > 1 - 2g + \sum k_j$. Moreover the coefficients $C_d(k, (a_j))$ form Vandermonde-type matrices and it can be proved that one can always find a set of positive integers $\{d_1, \cdots, d_l\}$ which will give linearly independent relations to solve for all the Hodge integrals of given dimension $3g - 2 + n$ with at-most one $\lambda$-class in terms of the values of lower-dim’l Hodge integrals with at-most one $\lambda$-class. So we just proved the following theorem;

**Theorem 7.2.** Any given Hodge integral with one $\lambda$-class:
\[
(14) \int_{\mathcal{M}_{g,n}} \psi_1^{k_1} \cdots \psi_n^{k_n} \lambda_j
\]
where $k_1, \cdots, k_n \in \mathbb{N} \cup \{0\}$, $j \in \{0, 1, 2, \cdots, g\}$, is explicitly expressed as a polynomial in terms of lower-dimensional Hodge integrals with one $\lambda$-class. Therefore it computes all Hodge integrals with one $\lambda$-class.

8. **Algorithm to compute Hodge Integrals with one $\lambda$-class**

It is clear that Theorem 7.2 can be implemented. We can use the formula (11) to compute Double Hurwitz numbers. And since the values of Hodge integrals will be too big to fit in the usual 4-byte integer data type, we will need a library for multi-precision computing. There are several free libraries on the web and GNU-MP is one
of them which provides well-organized C++ class interfaces as well as documentations.

For any given Hodge integral

\[ \int_{\mathcal{M}_{g,n}} \psi_1^{k_1} \cdots \psi_n^{k_n} \chi_j \]

Let \( e = (k_1, \ldots, k_{n-1}) \) and \( \chi = 2 - 2g \). Start with \( d = \chi - 1 + |e|, \cdots \) and find linear relations \( (10) \) as follows: Run over partitions \( \nu \) of size \( d \). For a fixed \( \nu \), run over pairs \( (\chi_0, \chi_\infty) \) which satisfies \( \chi = \chi_0 + \chi_\infty - 2l(\nu) \), \( \chi_\infty \leq 2 \), \( \chi_0, \chi_\infty \in 2\mathbb{Z} \). Now for a fixed pair \( (\chi_0, \chi_\infty) \), we can compute Double Hurwitz Number \( D^\bullet_{\infty}( (d), \nu) \) as follows: When \( c_1 \geq 0 \), i.e. when there are one or more 1’s in \( \nu \), we have

\[ [t^{2g}] \prod_{k \geq 1} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} = \sum \prod \sum 2^{2b_k} \frac{k^{2b_k}}{\prod_j (2a_j^k + 1)!} \]

where \( b_k = 0 \) if \( c_k = 0 \) and \( \chi_\infty = 2 - 2g_\infty \), \( \sum_k b_k = g_\infty \), \( \sum_j a_j^k = b_k \) with \( b_k, a_j^k \in \mathbb{N} \cup \{0\} \). When \( c_1 = -1 \), i.e. when there is no 1 in \( \nu \), let \( h \) be the minimum numbered part of \( \nu \). Then we have \( c_h \geq 1 \) and

\[ \left( \frac{\sinh(t/2)}{t/2} \right)^{-1} \left( \frac{\sinh(ht/2)}{ht/2} \right) = \frac{1}{h} (e^{(h-1)t/2} + e^{(h-3)t/2} + \cdots + e^{(-h+1)t/2}) \]

\[ = \left\{ \begin{array}{ll}
\sum_{m=0}^\infty \left[ \frac{1}{h^{2m-1}(2m)!} \left( \sum_l l^{2m} \right) \right] t^{2m}, & \text{when } h \text{ is even, } l = 1, 3, \cdots, h - 1 \\
\frac{1}{h} + \sum_{m=0}^\infty \left[ \frac{1}{h^{2m-1}(2m)!} \left( \sum_l l^{2m} \right) \right] t^{2m}, & \text{when } h \text{ is odd, } l = 2, 4, \cdots, h - 1
\end{array} \right. \]

Hence we have

\[ \prod_{k \geq 1} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} = \left( \frac{\sinh(t/2)}{t/2} \right)^{-1} \left( \frac{\sinh(ht/2)}{ht/2} \right) \left( \frac{\sinh(ht/2)}{ht/2} \right)^{c_h - 1} \prod_{k > h} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} \]

\[ = \left( \frac{1}{h} \delta_{h,\text{odd}} + \sum_{m=0}^\infty \frac{1}{h^{2m-1}(2m)!} \left( \sum_l l^{2m} \right) t^{2m} \right) \left( \frac{\sinh(ht/2)}{ht/2} \right)^{c_h - 1} \prod_{k > h} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} \]

and in this case, for \( g > 0 \);

\[ [t^{2g}] \prod_{k \geq 1} \left( \frac{\sinh(kt/2)}{kt/2} \right)^{c_k} = \sum (b_k) \left[ \frac{1}{h^{2b_k-1}(2b_k)!} \left( \sum_l l^{2b_k} \right) \right] \prod_{b_k \neq 0, k > h} 2^{2b_k} \frac{k^{2b_k}}{\prod (2a_j^k + 1)!} \]

where \( b_h = 0 \) if \( c_h = 1 \) and \( b_k = 0 \) if \( c_k = 0 \) for \( k \neq h \). Using these formulas and \( (11) \), we can compute Double Hurwitz numbers. In order to compute \( [\lambda^{(\nu)-\chi_0}] D^\bullet_{\nu,e}(\lambda) \), first run over the number of vertices \( m = 1, 2, \cdots, l(\nu) \). For each \( m \), find all possible groupings of \( \nu, e, \lambda \), and all possible splittings of \( \chi_0 \). Note that \( e \) admits groupings with empty components. Now find all triples \( (\nu(v), e(v), g(v)) \) according to the equivalence condition of vertices discussed in Section 3. Then the contribution of \( [\lambda^{(\nu)-\chi_0}] D^\bullet_{g(v),\nu(v),e(v)}(\lambda) \) will be the product of the combination factor \( 1/\prod m_i! \) and the expansions of \( D_{g(v),\nu(v),e(v)}(\lambda) \)
which can be obtained by

\[
\int_{\mathcal{M}_{g,l(\nu)+l(e)}} \frac{\Lambda^\nu g(1) \prod_{j=1}^{l(e)} (1 - \psi_j)^{e_j}}{\prod_{i=1}^{l(\nu)} (1 - \nu_i \psi_i)}
\]

\[
= \int_{\mathcal{M}_{g,l(\nu)+l(e)}} [1 - \lambda_1 + \cdots + (-1)^g \lambda_g] \prod_{j=1}^{l(e)} (\sum_{(l_j)} (-1)^{\tilde{l}_j} \frac{e_j}{\tilde{l}_j}) \prod_{i=1}^{l(\nu)} (\sum_{l_i} \nu_i \psi_i^{l_i})
\]

\[
= \sum_{k,(l_i),(\tilde{l}_j)} (-1)^{k+\sum \tilde{l}_j} \prod_{j=1}^{l(e)} (\frac{e_j}{\tilde{l}_j}) \prod_{i=1}^{l(\nu)} \nu_i \psi_i^{l_i} \int_{\mathcal{M}_{g,l(\nu)+l(e)}} \lambda_k \prod_{i=1}^{l(\nu)} \psi_i^{l_i} \prod_{j=1}^{l(e)} \psi_j^{\tilde{l}_j}
\]

where \(l_i \geq 0, 0 \leq \tilde{l}_j \leq e_j, 0 \leq k \leq g, \) and \(k + \sum_i l_i + \sum_j \tilde{l}_j = 3g - 3 + l(\nu) + l(e)\).

Some of them will have maximum dimension \(3g - 3 + n\) for the situations described in Section 7. Those are treated as unknowns and all others are lower-dimensional Hodge integrals or initial values which are already computed. Summing over all pairs \((\chi_0, \chi_\infty)\) and \(\nu\) will give a linear relation between Hodge integrals of the dimension \(3g - 3 + n\). Now we can follow same step as above for other values of \(d\) and obtain more linear relations. Observe that the number of unknowns are independent of \(d\) and actually bounded by the number of partitions of \(3g - 3 + n\), and hence we will have a system of linear relations which can be solved by simple Gaussian Elimination method. Thus in each dimension, it amounts to solve a matrix equation of size \(N \times N\) when \(N\) is at-worst-case the number of partitions of dimension.

9. Examples

In this section, I will show how the algorithm developed in the previous section works and check the results with the previously known-methods.

**Dimension 1:** There are 3 Hodge integrals with one \(\lambda\)-class of dimension 1;

\[
\int_{\mathcal{M}_{0,4}} \psi^1, \quad \int_{\mathcal{M}_{1,1}} \lambda_1, \quad \int_{\mathcal{M}_{1,1}} \psi^1
\]

And these integrals can be computed as follows:

\[
d = 3, g = 0, e = (0, 0, 0); \quad 6 \int_{\mathcal{M}_{0,4}} \psi^1 - 6 = 0 \implies \int_{\mathcal{M}_{0,4}} \psi^1 = 1
\]

\[
d = 1, g = 1, e = \emptyset, \quad \text{and} \quad d = 2, g = 1, e = \emptyset;
\]

\[
- \int_{\mathcal{M}_{1,1}} \lambda_1 + \int_{\mathcal{M}_{1,1}} \psi^1 = 0, \quad - \int_{\mathcal{M}_{1,1}} \lambda_1 + 3 \int_{\mathcal{M}_{1,1}} \psi^1 - \frac{1}{12} = 0
\]

\[
\implies \int_{\mathcal{M}_{1,1}} \lambda_1 = \int_{\mathcal{M}_{1,1}} \psi^1 = \frac{1}{24}
\]
This value of $\psi$ class integral recovers the exceptional case (2.46) in [21], and
the value of $\lambda_1$ class integral matches with the $\lambda_g$-formula since $B_2 = 1/6$.

**Dimension 2**: There are 5 Hodge integrals with one $\lambda$-class of dimension 2;

\[
\int_{\mathcal{M}_{0,5}} \psi^2, \quad \int_{\mathcal{M}_{0,5}} \psi^1 \psi^1, \quad \int_{\mathcal{M}_{1,2}} \psi^1 \lambda_1, \quad \int_{\mathcal{M}_{1,2}} \psi^2, \quad \int_{\mathcal{M}_{1,2}} \psi^1 \psi^1
\]

And these integrals can be computed as follows:

- $d = 3, g = 0, e = (0, 0, 0, 0)$;
  \[25 \int_{\mathcal{M}_{0,5}} \psi^2 - 25 = 0 \implies \int_{\mathcal{M}_{0,5}} \psi^2 = 1\]

- $d = 4, g = 0, e = (1, 0, 0, 0)$;
  \[-10 \int_{\mathcal{M}_{0,5}} \psi^1 \psi^1 + 20 = 0 \implies \int_{\mathcal{M}_{0,5}} \psi^1 \psi^1 = 2\]

- $d = 1, g = 1, e = (0)$, and $d = 2, g = 1, e = (0)$;
  \[-\int_{\mathcal{M}_{1,2}} \psi^1 \lambda_1 + \int_{\mathcal{M}_{1,2}} \psi^2 = 0, \quad -3 \int_{\mathcal{M}_{1,2}} \psi^1 \lambda_1 + 7 \int_{\mathcal{M}_{1,2}} \psi^2 - \frac{1}{6} = 0\]

  \[\implies \int_{\mathcal{M}_{1,2}} \psi^1 \lambda_1 = \int_{\mathcal{M}_{1,2}} \psi^2 = \frac{1}{24}\]

- $d = 2, g = 1, e = (1)$;
  \[-3 \int_{\mathcal{M}_{1,2}} \psi^1 \psi^1 + \frac{1}{8} = 0 \implies \int_{\mathcal{M}_{1,2}} \psi^1 \psi^1 = \frac{1}{24}\]

The first two values matches with the Witten's formula for the $g = 0$ case which says that

\[(15) \quad \int_{\mathcal{M}_{0,n}} \psi^{k_1} \cdots \psi^{k_n} = \binom{n - 3}{k_1, \ldots, k_n}, \quad \text{when } k_1 + \cdots + k_n = n - 3\]

and the remaining values match with the results in [31], and $\lambda_g$-formula (2).

**Dimension 3**: There are 8 Hodge integrals with one $\lambda$-class of dimension 3;

\[
\int_{\mathcal{M}_{0,6}} \psi^3, \quad \int_{\mathcal{M}_{0,6}} \psi^2 \psi^1, \quad \int_{\mathcal{M}_{0,6}} \psi^1 \psi^1 \psi^1, \quad \int_{\mathcal{M}_{0,6}} \psi^2 \lambda_1, \\
\int_{\mathcal{M}_{1,3}} \psi^3, \quad \int_{\mathcal{M}_{1,3}} \psi^1 \psi^1 \lambda_1, \quad \int_{\mathcal{M}_{1,3}} \psi^2 \psi^1, \quad \int_{\mathcal{M}_{1,3}} \psi^1 \psi^1 \psi^1
\]
And these integrals can be computed as follows:

\[ d = 3, g = 0, e = (0, 0, 0, 0, 0); \]
\[ 90 \int_{M_{0,6}} \psi^3 - 90 = 0 \implies \int_{M_{0,6}} \psi^3 = 1 \]

\[ d = 4, g = 0, e = (1, 0, 0, 0, 0); \]
\[ -65 \int_{M_{0,6}} \psi^2 \psi^1 + 195 = 0 \implies \int_{M_{0,6}} \psi^2 \psi^1 = 3 \]

\[ d = 5, g = 0, e = (1, 1, 0, 0, 0); \]
\[ 15 \int_{M_{0,6}} \psi^1 \psi^1 \psi^1 - 90 = 0 \implies \int_{M_{0,6}} \psi^1 \psi^1 \psi^1 = 6 \]

\[ d = 1, g = 1, e = (0, 0), \quad \text{and} \quad d = 2, g = 1, e = (0, 0); \]
\[ -\int_{M_{1,3}} \psi^2 \lambda_1 + \int_{M_{1,3}} \psi^3 = 0, \quad -7 \int_{M_{1,3}} \psi^2 \lambda_1 + 15 \int_{M_{1,3}} \psi^3 - \frac{1}{3} = 0 \]
\[ \implies \int_{M_{1,3}} \psi^2 \lambda_1 = \int_{M_{1,3}} \psi^3 = \frac{1}{24} \]

\[ d = 2, g = 1, e = (1, 0), \quad \text{and} \quad d = 3, g = 1, e = (1, 0); \]
\[ 3 \int_{M_{1,3}} \psi^1 \psi^1 \lambda_1 - 7 \int_{M_{1,3}} \psi^2 \psi^1 + \frac{1}{3} = 0, \]
\[ 6 \int_{M_{1,3}} \psi^1 \psi^1 \lambda_1 - 25 \int_{M_{1,3}} \psi^2 \psi^1 + \frac{19}{12} = 0 \]
\[ \implies \int_{M_{1,3}} \psi^1 \psi^1 \lambda_1 = \int_{M_{1,3}} \psi^2 \psi^1 = \frac{1}{12} \]

\[ d = 3, g = 1, e = (1, 1); \]
\[ 6 \int_{M_{1,3}} \psi^1 \psi^1 \psi^1 - \frac{1}{2} = 0 \implies \int_{M_{1,3}} \psi^1 \psi^1 \psi^1 = \frac{1}{12} \]

The first 3 values match with \(15\), and the rest match with \(\lambda_g\)-formula \(2\) and the results in \(31\), p36.
Dimension 4: There are 16 Hodge integrals with one \( \lambda \)-class of dimension 4;

\[
\begin{align*}
\int_{\mathcal{M}_{0,7}} \psi^4, & \quad \int_{\mathcal{M}_{0,7}} \psi^3 \psi^1, & \quad \int_{\mathcal{M}_{0,7}} \psi^2 \psi^2, & \quad \int_{\mathcal{M}_{0,7}} \psi^2 \psi^1 \psi^1, \\
\int_{\mathcal{M}_{4,1,4}} \psi^1 \psi^1 \psi^1, & \quad \int_{\mathcal{M}_{4,1,4}} \psi^3 \lambda_1, & \quad \int_{\mathcal{M}_{4,1,4}} \psi^4, & \quad \int_{\mathcal{M}_{4,1,4}} \psi^2 \psi^1 \lambda_1, \\
\int_{\mathcal{M}_{4,1,4}} \psi^3 \psi^1, & \quad \int_{\mathcal{M}_{4,1,4}} \psi^1 \psi^1 \lambda_1, & \quad \int_{\mathcal{M}_{4,1,4}} \psi^2 \psi^1 \psi^1, & \quad \int_{\mathcal{M}_{4,1,4}} \psi^2 \psi^2, \\
\int_{\mathcal{M}_{4,1,4}} \psi^1 \psi^1 \psi^1, & \quad \int_{\mathcal{M}_{2,1}} \psi^2 \lambda_2, & \quad \int_{\mathcal{M}_{2,1}} \psi^3 \lambda_1, & \quad \int_{\mathcal{M}_{2,1}} \psi^4.
\end{align*}
\]

And these integrals can be computed as follows:

\( d = 3, g = 0, e = (0, 0, 0, 0, 0, 0) \):

\[
301 \int_{\mathcal{M}_{0,7}} \psi^4 - 301 = 0 \implies \int_{\mathcal{M}_{0,7}} \psi^4 = 1
\]

\( d = 4, g = 0, e = (1, 0, 0, 0, 0, 0) \):

\[
-350 \int_{\mathcal{M}_{0,7}} \psi^3 \psi^1 + 1400 = 0 \implies \int_{\mathcal{M}_{0,7}} \psi^3 \psi^1 = 4
\]

\( d = 5, g = 0, e = (2, 0, 0, 0, 0, 0) \):

\[
140 \int_{\mathcal{M}_{0,7}} \psi^2 \psi^2 - 840 = 0 \implies \int_{\mathcal{M}_{0,7}} \psi^2 \psi^2 = 6
\]

\( d = 5, g = 0, e = (1, 1, 0, 0, 0, 0) \):

\[
140 \int_{\mathcal{M}_{0,7}} \psi^2 \psi^1 \psi^1 - 1680 = 0 \implies \int_{\mathcal{M}_{0,7}} \psi^2 \psi^1 \psi^1 = 12
\]

\( d = 6, g = 0, e = (1, 1, 1, 0, 0, 0) \):

\[
-21 \int_{\mathcal{M}_{0,7}} \psi^1 \psi^1 \psi^1 \psi^1 + 504 = 0 \implies \int_{\mathcal{M}_{0,7}} \psi^1 \psi^1 \psi^1 \psi^1 = 24
\]

These values match with (15).
These values match with [31], p36. and the $\lambda_g$-formula [2].
$d = 1, g = 2, e = \emptyset, \quad d = 2, g = 2, e = \emptyset, \quad$ and $d = 3, g = 2, e = \emptyset$; 

$$\int_{\mathcal{M}_{g,1}} \psi^2 \lambda_2 - \int_{\mathcal{M}_{g,1}} \psi^3 \lambda_1 + \int_{\mathcal{M}_{g,1}} \psi^4 = 0,$$

$$7 \int_{\mathcal{M}_{g,1}} \psi^2 \lambda_2 - 15 \int_{\mathcal{M}_{g,1}} \psi^3 \lambda_1 + 31 \int_{\mathcal{M}_{g,1}} \psi^4 - \frac{1}{240} = 0,$$

$$25 \int_{\mathcal{M}_{g,1}} \psi^2 \lambda_2 - 90 \int_{\mathcal{M}_{g,1}} \psi^3 \lambda_1 + 301 \int_{\mathcal{M}_{g,1}} \psi^4 - \frac{5}{48} = 0$$

$$\Rightarrow \int_{\mathcal{M}_{g,1}} \psi^2 \lambda_2 = \frac{7}{5760}, \quad \int_{\mathcal{M}_{g,1}} \psi^3 \lambda_1 = \frac{1}{480}, \quad \int_{\mathcal{M}_{g,1}} \psi^4 = \frac{1}{1152}$$

The first value matches with $\lambda_g$-formula (2), since we have

$$\sum_{k=0}^{m} \binom{m+1}{k} B_k = 0, \quad \text{for} \ m > 0 \Rightarrow \ B_4 = -\frac{1}{30}$$

$$\left(\frac{2 \times 2 + 1 - 3}{2} \right) \frac{2^{2g-1} - 1}{2^{2g-1}} \frac{|B_{2g}|}{(2 \times 2)!} = \frac{7}{5760}$$

The second value matches with (3), since we have

$$b_g = \frac{2^{2g-1} - 1}{2^{2g-1}(2g)!} |B_{2g}|, \quad \text{for} \ g > 0 \Rightarrow b_1 = \frac{1}{24}, \quad b_2 = \frac{7}{5760}$$

$$b_g \sum_{i=1}^{2g-1} \frac{1}{i} - \frac{1}{2} \sum_{g_1+g_2=g} \frac{(2g_1-1)! (2g_2-1)!}{(2g-1)!} b_{g_1} b_{g_2} = \frac{7}{5760} \left(1 + \frac{1}{2} + \frac{1}{3}\right) - \frac{1}{2} \frac{1}{3} \frac{1}{24} = \frac{1}{480}$$

And the last value matches with the result in [31], p36.

REFERENCES

[1] I.P. Goulden, D.M. Jackson, A. Vainshtein, The number of ramified coverings of the sphere by the torus and surfaces of higher genera, Ann. of Comb. 4 (2000), 27-46.
[2] I.P. Goulden, D.M. Jackson, R. Vakil, Towards the Geometry of Double Hurwitz Numbers, preprint, math.AG/0309440
[3] E. Getzler, Topological recursion relations in genus 2, in “Integrable systems and algebraic geometry” (Kobe/Kyoto, 1997), World Science Publishing, River Edge, NJ. 1998, pp 73-106
[4] C. Faber, Algorithms for computing intersection numbers on moduli spaces of curves, with an application to the class of the locus of Jacobians, in New Trends in Algebraic Geometry (K.Hulek, F.Catanese, C.Peters and M.Reid, eds.), 93-109, Cambridge University Press, 1999
[5] C. Faber, R. Pandharipande, Hodge integrals, partition matrices, and the $\lambda_g$ conjecture, Ann. of Math. (2) 157 (2003), no. 1, 97–124
[6] T. Graber, R. Pandharipande, Localization of virtual classes, Invent. Math. 135 (1999), no. 2, 487–518.
[7] T. Graber, R. Vakil, Relative virtual localization and vanishing of tautological classes on moduli spaces of curves, preprint, math.AG/0309227
[8] E.-N. Ionel, T. Parker, *Relative Gromov-Witten invariants*, Ann. of Math. (2) **157** (2003), no. 1, 45–96.

[9] E.-N. Ionel, T. Parker, *The symplectic sum formula for Gromov-Witten invariants*, preprint, math.SG/0010207.

[10] M. Kontsevich, *Intersection theory on the moduli space of curves and the matrix Airy function*, Comm. Math. Phys. **147** (1992), no. 1, 1–23.

[11] J. Li, *Stable Morphisms to singular schemes and relative stable morphisms*, J. Diff. Geom. **57** (2001), 509-578.

[12] J. Li, *Relative Gromov-Witten invariants and a degeneration formula of Gromov-Witten invariants*, J. Diff. Geom. **60** (2002), 199-293.

[13] A.M. Li, G. Zhao, Q. Zheng, *The number of ramified coverings of a Riemann surface by Riemann surface*, Comm. Math. Phys. **213** (2000), no. 3, 685–696.

[14] C.-C. Liu, K. Liu, J. Zhou, *A proof of a conjecture of Mariño-Vafa on Hodge Integrals*, J. Differential Geom. **65** (2003), no. 2, 289–340.

[15] C.-C. Liu, K. Liu, J. Zhou, *A formula of two-partition Hodge integrals*, preprint, math.AG/0310272.

[16] C.-C. Liu, K. Liu, J. Zhou, *Mariño-Vafa formula and Hodge Integral Identities*, preprint, math.AG/0308015.

[17] I.G. MacDonald, *Symmetric functions and Hall polynomials*, 2nd edition. Claredon Press, 1995.

[18] M. Mariño, C. Vafa, *Framed knots at large N*, Orbifolds in mathematics and physics (Madison, WI, 2001), 185–204, Contemp. Math., 310, Amer. Math. Soc., Providence, RI, 2002.

[19] A. Okounkov, R. Pandharipande, *Hodge integrals and invariants of the unknot*, Geom. Topol. **8** (2004), 675–699.

[20] E. Witten, *Quantum field theory and the Jones polynomial*, Commun. Math. Phys. **121** (1989) 351–399.

[21] E. Witten, *Two-dimensional gravity and intersection theory on moduli space*, Surveys in differential geometry (Cambridge, MA, 1990), 243–310, Lehigh Univ., Bethlehem, PA, 1991.

[22] J. Zhou, *Hodge integrals, Hurwitz numbers, and symmetric groups*, preprint, math.AG/0308024.

[23] J. Zhou, *A conjecture on Hodge integrals*, preprint, math.AG/0310282.

[24] J. Zhou, *Localizations on moduli spaces and free field realizations of Feynman rules*, preprint, math.AG/0310283.

[25] J. Zhou, *Hodge Integrals and Integrable Hierarchies*, preprint, math.AG/0310408.

[26] M. Mulase, *Complete Integrability of the Kadomtsev-Petviashvili Equation*, Adv. Math. **54** (1985), 57-66.

[27] M. Mulase, *Cohomological Structure in Soliton Equations and Jacobian Varieties*, J. Diff. Geom. **19** (1984), 403-430.

[28] I.B. Frenkel, W. Wang, *Virasoro algebra and wreath product convolution*, J. Alg. **242** (2001), 656-671.

[29] T. Miwa, M. Jimbo, E. Date, *Solitons: Idifferential equations, symmetries and infinite-dimensional algebras*, Cambridge Tracs in Mathematics, 135, Cambridge University Press, 2000.

[30] M. Mirzakhani, *Simple geodesics and Weil-Petersson volumes of moduli spaces of bordered Riemann surfaces*, preprint, 2003.

[31] C. Itzykson, J.-B. Zuber, *Combinatorics of the Modular Group II: The Kontsevich integrals*, preprint, hep-th/9201001.

University of California, Los Angeles, CA 90095-1555

E-mail address: yskin@math.ucla.edu