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Detection of faint broad emission lines in type 2 AGN – II. On the measurement of the black hole mass of type 2 AGN and the unified model

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ABSTRACT

We report the virial measurements of the black hole (BH) mass of a sample of 17 type 2 active galactic nuclei (AGN), drawn from the Swift/BAT 70-month 14–195 keV hard X-ray catalogue, where a faint BLR component has been measured via deep NIR (0.8–2.5 μm) spectroscopy. We compared the type 2 AGN with a control sample of 33 type 1 AGN. We find that the type 2 AGN BH masses span the 5 < log(MBH/M☉) < 7.5 range, with an average log(MBH/M☉) = 6.7, which is ~0.8 dex smaller than found for type 1 AGN. If type 1 and type 2 AGN of the same X-ray luminosity log(L14–195/erg s⁻¹) ~ 43.5 are compared, type 2 AGN have 0.5 dex smaller BH masses than type 1 AGN. Although based on few tens of objects, this result disagrees with the standard AGN unification scenarios in which type 1 and type 2 AGN are the same objects observed along different viewing angles with respect to a toroidal absorbing material.

Key words: galaxies: active – quasars: emission lines – quasars: supermassive black holes – galaxies: Seyfert – infrared: galaxies.

1 INTRODUCTION

In the last decade, by using hard X-ray (2-10 keV) selected active galactic nuclei (AGN) samples, it has been possible to accurately derive the AGN luminosity function up to z ~ 6 (e.g. Ueda et al. 2014). Moreover, by using virial-based techniques in the optical band on samples of broad line type 1 AGN (AGN1), it has been possible to estimate the supermassive black hole (BH) mass function (e.g. Greene & Ho 2007; Kelly et al. 2010; Schulze et al. 2015). However, the BH mass (MBH) measurements are affected by several selection biases against the narrow line type 2 AGN (AGN2), or low-luminosity AGN, where the broad-line region (BLR) is not visible in the rest-frame optical band because of either dust absorption or dilution by the host galaxy spectra (see e.g. Baldassare et al. 2016).

According to the original standard unified model (Antonucci 1993), the different observational classes of AGN (AGN1 and AGN2) are believed to be the same kind of objects observed under different conditions (e.g. different orientations of the observer with respect to a dusty torus). In the framework of the AGN phenomenon and co-evolution, this implies that AGN with the same luminosity should share, on average, the same properties (e.g. same masses, same accretion rates and, then, same Eddington ratios λEdd = Lbol/LEdd).

Nevertheless, nowadays there is growing evidence that AGN1 and AGN2 could belong to intrinsically different populations (see e.g. Elitzur 2012; Lanzuisi et al. 2015), having, on average, different luminosities (lower for AGN2; Ueda et al. 2003; La Franca et al. 2005; Ueda et al. 2014), different accretion rates (smaller for AGN2, Winter et al. 2010; Lusso et al. 2012), different host galaxy properties (more late-type for AGN2), and different clustering, environment and halo mass properties (Allevato et al. 2014). The observed difference in the luminosity distributions of AGN1 and AGN2 could, however, still comply with an orientation-based unified model in which the torus opening angle (or the absorbing material covering factor) depends on luminosity. In this scenario, many of the observed differences between the AGN1 and the AGN2...
population can be attributed to selection effects. In contrast, if a difference is measured in the average $M_{BH}$ (or host halo mass and clustering properties) of AGN1 and AGN2 sharing the same intrinsic (corrected for absorption) luminosity, then AGN1 and AGN2 should be intrinsically different objects and the unified model should be revised.

We have therefore started a project aimed at measuring the BH mass in AGN2 (La Franca et al. 2015, 2016; Onori et al. 2017; Ricci et al. 2017). In those few studies where AGN2 BH masses have been derived (e.g. Heckman et al. 2004, from The Sloan Digital Sky Survey), the authors used the $M_{BH}$-host scaling relations that have not yet been hold to provide for such a population (see Graham 2008; Kormendy, Bender & Cornell 2011). Several studies have shown that many AGN2 exhibit faint broad-line components if observed with high ($\geq 20)$ S/N in the near-infrared (NIR; 0.8–2.5 $\mu$m), where the dust absorption is less severe than in the optical (Veilleux, Goodrich & Hill 1997; Riffel, Rodríguez-Ardila & Pastoriza 2006; Cai et al. 2010). Moreover, observation in the NIR of AGN1, whose $M_{BH}$ were measured using reverberation mapping techniques (RM; Blandford & McKee 1982; Peterson 1993) have demonstrated that the virial method to measure $M_{BH}$ can be efficiently used with the NIR Pa $\alpha$ $\lambda$1.875 $\mu$m and Pa$\beta$ $\lambda$1.282 $\mu$m lines (Kim, Im & Kim 2010; Landt et al. 2011).

Following the above studies and, as it is observed for the optical continuum luminosity, the empirical relation between the X-ray luminosity and the dimension of the BLR ($R_{BLR} \propto L_{X}$; Maiolino et al. 2007; Greene et al. 2010, Ricci et al. 2017) have calibrated new virial relations between the FWHM$_{NIR}$ of the most relevant NIR emission lines (Pa $\alpha$, Pa$\beta$, He $\i$ $\lambda$1.083 $\mu$m) and the intrinsic hard X-ray luminosity, $L_{X}$, of the type $M_{BH} \propto$ FWHM$_{NIR}$ $\propto L_{X}^{0.5}$. These relations can be used to measure $M_{BH}$ of either AGN2 or obscured and low-luminosity AGN1 (see also La Franca et al. 2015, 2016).

In this paper, we present the measurement for $M_{BH}$ of a sample of AGN2 selected from the Swift/BAT 70-month catalogue ( Baumgartner et al. 2013), where a faint BLR component in the NIR emission line was found (Onori et al. 2017). We compare the resulting $M_{BH}$ distribution with that of a control sample of AGN1 selected from the Swift/BAT 70-month catalogue and whose $M_{BH}$ have been measured via RM techniques. We adopt a $\Delta_{m} = 0.3$, $\Omega_{m} = 0.7$ and $H_{0} = 70$ km s$^{-1}$ Mpc$^{-1}$ cosmology. Unless otherwise stated, all the quoted uncertainties are at 68 per cent confidence level.

2 DATA AND SAMPLES
In order to measure the BH mass of the AGN2, we have carried out NIR spectroscopic observations of 41 obscured and intermediate-class AGN (type 2, 1.9 and 1.8; all named AGN2 in the following) at redshift $z \lesssim 0.1$, randomly selected from the Swift/BAT 70-month catalogue. Thanks to the very hard X-ray band (14–195 keV) that has been used to build the catalogue, the parent AGN sample is free of absorption selection effects up to log$N_{HI} \leq$ 24 cm$^{-2}$ (see e.g. fig. 11 in Burlon et al. 2011). The observations have been carried out in the framework of a systematic study of the AGN2 NIR spectral properties and have been executed using ISAAC/VLT, X-shooter/VLT and LUCI/LBT spectrometers, with a spectral resolution of 65, 20–78 and 220 km s$^{-1}$, respectively, and reaching an average S/N ratio of $\geq$ 30 per resolution element. A BLR component showing 800 $\lesssim$ FWHM $\lesssim$ 3500 km s$^{-1}$, significantly larger than the narrow-line region component measured in the same spectra, has been identified in 13 out of 41 AGN2 ($\sim$30 per cent of the total sample). The data reduction, spectral analysis and line fitting parameters (full width at half–maximum, FWHM, and fluxes of the most relevant emission lines) have been published in a companion paper (Onori et al. 2017). The sample has been extended with four AGN2 included in the Swift/BAT 70-month catalogue, whose FWHM NIR lines, or spectra, were available in the literature. The spectral data, when available, have been fitted using the technique described in Onori et al. (2017). In Table 1, we list the FWHM of the BLR components of the He $\i$ and Pa$\beta$ lines of the AGN2 sample (in the following called AGN2/BLR sample).

In Fig. 1 (top panel), we show the $L_{14–195}$ distributions of all the AGN2 included in the Swift/BAT 70-month catalogue (red dotted–dashed line) as well as the distributions of the 17 AGN2/BLR (red continuous line). The average X-ray luminosity is log$L_{14–195} = 43.5$ with a spread $\sigma$ (log$L_{14–195}$) = 0.9 and log$L_{14–195} = 42.9$ erg s$^{-1}$ with a spread $\sigma$ (log$L_{14–195}$) = 0.9, for the Swift/BAT AGN2 and the AGN2/BLR, respectively.

In order to build a control sample of AGN1 to be compared with the results obtained from the AGN2 population, we have used 33 AGN1 included in the Swift/BAT 70-month catalogue and whose $M_{BH}$ have been measured via RM techniques. This sample includes those 31 AGN1 selected by Ricci et al. (2017, Table 1) plus two additional RM AGN1, namely 3C 390.3 (log$L_{14–195} = 44.88$ erg s$^{-1}$, $M_{BH} = 278^{+24}_{-32} \times 10^{8}$ M$\odot$) and Mrk 50 (log$L_{14–195} = 43.45$ erg s$^{-1}$, $M_{BH} = 260^{+23}_{-32} \times 10^{7}$ M$\odot$).
Table 1. The sample of AGN2/BLR and their BH masses.

| Object name | $z$  | Class | $A_V$ (mag) | $W_2–W_4$ (mag) | log($L_{14–195}$) (erg s$^{-1}$) | He I (km s$^{-1}$) | Paβ (km s$^{-1}$) | NIR (km s$^{-1}$) | log($M_{BH}$) (M$_\odot$) | log($L_{bol}$) (erg s$^{-1}$) | $\lambda_{Edd}$ |
|-------------|------|-------|------------|----------------|-----------------|----------------|-------------|----------------|----------------|----------------|-------------|
| 2MASX J05054575–2351139 | 0.0350 | 2 | – | 5.80 | 44.24 | 1772$^{+319}_{-232}$ | ... | ... | 1772$^{+319}_{-232}$ | 7.37 ± 0.18 | 45.50 ± 0.849 |
| 2MASX J18305065+0928414 | 0.0190 | 2 | – | 4.45 | 42.40 | 3513$^{+213}_{-232}$ | ... | ... | 3513$^{+213}_{-232}$ | 7.04 ± 0.09 | 43.09 ± 0.009 |
| ESO 234-G-050 | 0.0088 | 2 | 0.8 | 6.30 | 42.29 | 1110$^{+63}_{-59}$ | 1304$^{+381}_{-322}$ | 1128$^{+106}_{-106}$ | 6.00 ± 0.10 | 42.98 ± 0.076 |
| ESO 374-G-044 | 0.0284 | 2 | – | 6.96 | 43.57 | 1123$^{+383}_{-221}$ | 1412$^{+318}_{-294}$ | 1265$^{+215}_{-215}$ | 6.74 ± 0.15 | 44.71 ± 0.742 |
| MCG -01-24-12 | 0.0196 | 2 | – | 6.32 | 43.55 | – | 2069$^{+300}_{-280}$ | 2069$^{+300}_{-280}$ | 7.16 ± 0.12 | 44.24 ± 0.096 |
| MCG -05-23-16 | 0.0085 | 2 | 6.8 | 5.84 | 43.51 | 2474$^{+667}_{-667}$ | 2133$^{+93}_{-93}$ | 2278$^{+162}_{-162}$ | 7.22 ± 0.06 | 44.20 ± 0.075 |
| NGC 1052 | 0.0050 | 2 | 1.5$^a$ | 5.12 | 42.22 | 2417$^{+143}_{-128}$ | – | 2417$^{+143}_{-128}$ | 6.63 ± 0.09 | 42.91 ± 0.015 |
| NGC 1365 | 0.0055 | 1.8 | 5.2 | 7.31 | 42.63 | – | 1971$^{+85}_{-75}$ | 1971$^{+85}_{-75}$ | 6.65 ± 0.09 | 43.32 ± 0.037 |
| NGC 2992 | 0.0077 | 2 | 5.1$^a$ | 6.25 | 42.55 | 3157$^{+586}_{-400}$ | 2055$^{+20}_{-20}$ | 2218$^{+190}_{-190}$ | 6.72 ± 0.08 | 43.24 ± 0.026 |
| NGC 4395 | 0.0013 | 1.9 | 4.0 | 6.32 | 40.79 | 1332$^{+393}_{-110}$ | 851$^{+29}_{-34}$ | 990$^{+72}_{-72}$ | 5.14 ± 0.07 | 41.48 ± 0.017 |
| NGC 6221 | 0.0050 | 2 | 3.2 | 7.41 | 42.05 | 2141$^{+141}_{-128}$ | 2250$^{+80}_{-80}$ | 2195$^{+155}_{-155}$ | 6.46 ± 0.06 | 42.74 ± 0.015 |
| NGC 7314 | 0.0048 | 1.9 | 4.4 | 5.70 | 42.42 | 1427$^{+46}_{-38}$ | 1347$^{+16}_{-8}$ | 1384$^{+99}_{-99}$ | 6.24 ± 0.06 | 43.11 ± 0.058 |

Notes: (1) AGN name; (2)–(3) redshift, source classification (from Baumgartner et al. 2013); (4) optical–NIR extinction; (5) mid-IR 4.6–22 μm colour; (6) 14–195 keV intrinsic luminosity (from Baumgartner et al. 2013); (7)–(8) intrinsic (corrected for instrumental resolution) FWHM of the broad emission-line component of Hε and Paβ from Onori et al. (2017); (9) Average FWHM of the HeI and Paβ lines; (10) $M_{BH}$ derived with the virial relation a3 of table 4 from Ricci et al. (2017). A constant virial factor $f = 4.31$ (Grier et al. 2013) has been adopted. An uncertainty $\epsilon = 0.5$ dex, due to the spread of the population, should also be taken into account (Ricci et al. 2017). (11)–(12) bolometric luminosity and corresponding Eddington ratio, both calculated adopting the $K$-correction of Vasudevan & Fabian (2007).

$^a$From Bertuscher et al. (2015);
$^b$Measure of the FWHM of the Pa α line from Cai et al. (2010);
$^c$The FWHM have been measured applying the fitting procedure described by Onori et al. (2017) on the data published by Riffel et al. (2006);
$^d$From Ramos Almeida, Pérez García & Acosta-Pulido (2009).

In the luminosity range where the two distributions overlap, 42.5 < log($L_{14–195}$) < 44.5, AGN1 show significantly larger FWHM than AGN2: ~3400 km s$^{-1}$ instead of ~1970 km s$^{-1}$ (log(FWHM) = 3.53 ± 0.036 and log(FWHM) = 3.294 ± 0.032 for AGN1 and AGN2/BLR, respectively).

In order to compute the AGN2 $M_{BH}$, we have used the relation

$$\log \left( \frac{M_{BH}}{M_\odot} \right) = 7.75 + 0.5 \log \left( \frac{FWHM_{NIR}}{10^6 \text{ km s}^{-1}} \right)^2 \left( \frac{L_{14–195} \text{ keV}}{10^{42} \text{ erg s}^{-1}} \right)^{0.5},$$

The FWHM have been measured by Onori et al. (2017) on the data published by Riffel et al. (2006); from Ramos Almeida, Pérez García & Acosta-Pulido (2009).

Measurements of the BH mass of AGN2

3 COMPARISON BETWEEN THE AGN1 AND AGN2 POPULATIONS

As discussed by Landt et al. (2008), Onori et al. (2017) and Ricci et al. (2017), the most relevant NIR emission lines of the BLR (Pa α, Paβ, He I λ1.083 μm) have, within the errors, the same FWHM. Therefore, the average width (FWHM$_{ave}$) of these lines (when available). In Fig. 1 (top panel), we show the $L_{14–195}$ distribution of all the AGN1 included in the Swift/BAT 70-month catalogue (blue dotted line) as well as the distribution of the control sample of 33 AGN1 (blue dotted line). The average X-ray luminosity is log($L_{14–195}$) $\sim$ 44.3 with a spread $\sigma$ (log($L_{14–195}$)) = 1.0 and log($L_{14–195}$) $\sim$ 43.8 with a spread $\sigma$ (log($L_{14–195}$)) = 1.0, for the Swift/BAT AGN1 and the control sample of AGN1, respectively.
As also shown in Fig. 3, the $M_{\rm BH}$ of AGN2/BLR (red continuous line) is typically smaller than AGN1 (blue dotted line). The average (largest) $M_{\rm BH}$ is log($M_{\rm BH}/M_\odot$) $\sim$ 6.8 ($\sim$ 7.5) in the AGN2 sample, while in the control AGN1 sample, it is log($M_{\rm BH}/M_\odot$) $\sim$ 7.6 ($\sim$ 9.2). In the 42.5 < log($L_{14.195}$/erg s$^{-1}$) < 44.5 luminosity bin, the average $M_{\rm BH}$ of the AGN2 sample is ~0.5 dex smaller than that measured in the AGN1 sample (log($M_{\rm BH}/M_\odot$) = 7.08 $\pm$ 0.10 and log($M_{\rm BH}/M_\odot$) = 7.61 $\pm$ 0.01, for AGN2/BLR and AGN1, respectively). For the sub-sample of 20 AGN1 for which the FWHM$_{\rm BLR}$ measurements are available, the average $M_{\rm BH}$ obtained using equation (1), instead of using the RM measurements, is very similar: log($M_{\rm BH}/M_\odot$) = 7.63 $\pm$ 0.01.¹

The above result is also illustrated in Fig. 2 (bottom panel), where the $L_{14.195}/M_{\rm BH}$ ratio (which is a proxy of the Eddington ratio $\lambda_{\rm Edd}$), plus a constant, is shown as a function of $L_{14.195}$. We computed the Eddington ratio assuming the bolometric correction of Vasudevan & Fabian (2007)² (see Table 1). In the overlapping luminosity bin 42.5 < log($L_{14.195}$/erg s$^{-1}$) < 44.5, AGN2 have, on average, 0.3 dex larger $\lambda_{\rm Edd}$ than AGN1 (log$\lambda_{\rm Edd}$ $\simeq$ −0.85 and 1.15, for AGN2/BLR and AGN1, respectively).

4 ANALYSIS OF POSSIBLE SELECTION EFFECTS

Although based on few tens of objects, at face value, our results imply that AGN2 have, on average, about 0.25 dex lower BLR FWHM, and 0.5 dex lower $M_{\rm BH}$ (larger $\lambda_{\rm Edd}$) than AGN1 of the same luminosity. In Onori et al. (2017), we have investigated whether our FWHM measurements could be affected by some selection biases. No dependence was found in the sample where the BLR was measured on both the X-ray and NIR fluxes, on the orientation angle of the host galaxy or on the hydrogen column density, $N_\text{H}$, measured in the X-ray band.

A possible selection could be originated by the effects of the absorption/reddening medium along the line of sight, which is obviously present as it is at the origin of the AGN2 classification, and then large $N_\text{H}$ values (>10$^{21}$ cm$^{-2}$ s$^{-1}$) are typically measured in AGN2. One scenario could be that the most central parts of the BLR are embedded in a region of absorbing material and the broad components that we have detected originate in the outer, and therefore slower, part of the BLR. In this case, a trend should be visible where the largest FWHM are detected in the less X-ray absorbed and/or less reddened objects. As shown in Fig. 4, where we plot the FWHM as a function of $N_\text{H}$ and as a function of the extinction $A_V$, we do not find such trends in our AGN2 sample. The extinction $A_V$ has been estimated in eight AGN2 using either the BLR Paschen and Balmer line ratios (when available), assuming a Milky Way reddening law (Allen 1976) and $R_V$ = 3.1, or the values derived by Burtscher et al. (2015) using a ‘dust colour’ method (see Table 1). The measured $A_V$–$N_\text{H}$ distribution of our AGN2 sample is typical of the AGN2 population (see e.g. Burtscher et al. 2016, fig. 3).

¹ Indeed, equation (1) has been calibrated by Ricci et al. (2017) using almost the same RM AGN1 of the control sample (31 out of 37 objects are in common).

² As parametrized by Shankar, Weinberg & Miralda-Escudéro (2013, equation 22, in which, because of a typo, log$L_\text{x}$ should read log$L_{\text{bol}}$). No significant difference in the results was found by using the bolometric correction of Marconi et al. (2004).
The presence of dust inside the BLR is unlikely as it would disagree with the results obtained through dust RM at K-band wavelength, so that the time lag of the inner radius of the torus is three to four times longer than the \( \text{H}\beta \) time lag (Burlon et al. 2013; Vazquez et al. 2015, and references therein). The dust, instead, is consistent with being originated at the mass of the predicted graphite sublimation radius (Netzer & Laor 1993; Netzer 2015). This implies that the BLR is indeed bound by the dust distribution, as also observationally confirmed by Landt et al. (2014, but see Czerny & Hryniewicz 2011 for a different scenario in which the BLR is originated in regions where dust could co-exist).

We can therefore conclude that, once a BLR component is detected, the dusty region (maybe clumpy, according to recent studies; Marinucci et al. 2016, and references therein) that surrounds the AGN should have been completely penetrated. Note that a possible consequence of dust absorption could be the reduction of the line intensity (but not of the BLR FWHM; see also Kim et al. 2015a). This could affect the estimation of the BH mass, if the line luminosity were to be used, as it is the case of some virial relations for type 1 AGN. Therefore, in the case of obscured AGN, instead of the optical/NIR line luminosities, it is better to use the hard X-ray luminosities as a proxy of the BLR radius in the virial relations for the estimation of the BH mass.

Another possible bias could be due to the fact that in the standard unification model, obscured AGN are viewed, on average, at larger angles relative to the accretion disc axis than unobscured AGN. Therefore, the estimation of the BLR FWHM could be affected by projection effects. Indeed, the dependence of BLR linewidths in terms of orientations is observationally well known (Wills & Browne 1986). Recently, Bisogni, Marconi & Risaliti (2017), using the EW of the \([\text{O}\,\text{iii}] \lambda 5007 \, \text{Å} \) line as an inclination indicator, have found that the more inclined AGN1 have, on average, larger FWHM of the BLR \( \text{H}\beta \) line. These results match with those of Pancoast et al. (2014), who found that the virial \( f \) factor decreases with increasing inclination (see also Risaliti, Salvati & Marconi 2011). A smaller \( f \) factor is needed to compensate the increasing broadening of the FWHM with the inclination of the observed emission-line width. These results have been interpreted as a hint at a possible disc-like shape for the BLR (Bisogni et al. 2017).

However, the EW of the \([\text{O}\,\text{iii}] \) line is not a good indicator of the orientation for AGN2, as the AGN continuum component is suppressed and overwhelmed by the host galaxy contribution. Therefore, following Rose et al. (2015), we have used the \( W_{2–W4} \) (4.6–22 \( \mu \)m) mid-IR colour, as measured by the Wide-Field Infrared Survey Explorer (Wright et al. 2010), to roughly estimate the orientation of our samples of AGN. In Table 1, the \( W_{2–W4} \) colours of the AGN2/BLR sample are listed. The more inclined AGN2 should show redder \( W_{2–W4} \) colours than AGN1, because the hottest dust emission of the inner regions of the torus should be less visible as the inclination increases (Rose et al. 2015, and references therein). Our AGN2/BLR show a red \( W_{2–W4} \) distribution, typical of the AGN2 population, having, an average, \( < W_{2–W4} > = 6.2 \) with a \( (1\sigma) \) spread of 0.8, while the AGN1 control sample has \( < W_{2–W4} > = 5.2 \) with a \( (1\sigma) \) spread of 0.7 (see Fig. 5 and, for comparison, fig. 2 in Rose et al. 2015). If the common luminosity bin 42.5 < \( \log L_{14–195} \) erg s\(^{-1} \) < 44.5 is considered, AGN2/BLR (AGN1) have \( < W_{2–W4} > = 6.4 \) (5.4) with a \( (1\sigma) \) spread of 0.6 (0.7).

We can conclude that the smaller BLR FWHM and lower \( M_{\text{BH}} \) of our sample of AGN2 with respect to the AGN1 of the same luminosity should not be ascribed to orientation effects. Indeed, the larger inclination of the AGN2 should cause larger FWHM to be observed, and even smaller \( f \) factors to be used. Therefore, following Pancoast et al. (2014), even smaller \( M_{\text{BH}} \) could be derived for the AGN2/BLR sample than those obtained in this work, where a common \( f = 4.31 \) factor, for both AGN1 and AGN2, is assumed.

5 CONCLUSIONS

Determining the distribution of \( M_{\text{BH}} \) of AGN is of paramount importance in order to understand the AGN phenomenon.

Using deep NIR spectroscopic observations, we have detected faint BLR components in a sample of 17 AGN2 (Onori et al. 2017) drawn from the 14–195 keV X-ray \( \text{Swift}/\text{BAT} \) 70-month catalogue, which is free of biases against absorbed sources up to \( \log N_{\text{HI}} \lesssim 24 \, \text{cm}^{-2} \) (Burlon et al. 2011). No dependence was found on the BLR detection success rate, or FWHM of the lines, on either the X-ray or NIR fluxes, or on orientation angle of the host galaxy, on the hydrogen column density and on the extinction.

In this work, we have found that the average AGN2 FWHM of the BLR is \( \sim 0.25 \) dex smaller than measured in a control sample of AGN1 having the same average X-ray intrinsic luminosity. Using new virial relations calibrated by Ricci et al. (2017), which are based on the FWHM of the most relevant BLR emission lines and the intrinsic hard X-ray luminosity, we have measured the \( M_{\text{BH}} \) of the AGN2 in our sample. The \( M_{\text{BH}} \) of the AGN2 are, on average, \( \sim 0.8 \)
Our findings could instead fit in an evolutionary scenario (see e.g. Hopkins et al. 2005) in which AGN2 represents the preceding stage of a type 1 AGN. In this picture, AGN2 are dust-enshrouded ‘buried’, low-mass ($M_{BH} \lesssim 10^{-3} M_\odot$) BHs that accrete at high Eddington ratios. When feedback from the accretion drives away the obscuring material, a window is created in which the AGN is seen as an optical type 1. In this evolutionary path, highly accreting red quasars, as those observed by Kim et al. (2015b), could belong to an intermediate population. Eventually, the activity ends when the accretion rate drops below that required to maintain the typical AGN luminosities.

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