Pileup measurement and mitigation techniques in CMS

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Abstract. When trying to reconstruct an event from a hard-scatter pp collision in CMS, it is of the utmost importance to correctly measure the energy from jets. The jet energy corrections (JEC) correct, on average, the energy of the reconstructed jets back to the energy of the final-state particles that initiated the jets. This effort is hindered by additional energy in the jets coming from other soft pp collisions. The additional energy is termed pileup or offset and comes from everything except the primary vertex (PV) and its underlying event (UE). In this paper, we describe how this pileup energy is measured and parametrized as well as the techniques used to remove this extra energy from the reconstructed jets.

1. Introduction
During the reconstruction of a proton-proton (pp) collision (event), jets are often reconstructed with a $p_T$ that differs from that of the final-state particles that make up the jet. The jet energy corrections (JEC) correct the reconstructed jet energy back to the true energy of the final-state particles. All analyses at CMS which make use of jets must use the JEC. Therefore, the CMS collaboration has developed a factorized approach to these JEC which is composed of multiple levels representing corrections for various physics or detector effects and allows enough flexibility in the corrections to be applicable to many types of analyses [1–3].

One of the physics effects that the JEC seek to correct for is the additional energy attributed to the jets which comes from pp interactions other than the hard-scatter event at the primary vertex (PV). The additional energy from these events is often called pileup or offset and will be the focus of this paper.
Figure 2. A basic graphic showing two jets coming from a PV with additional $pp$ interactions in red. The arrows coming from the red $pp$ interactions are particles which add additional energy to the jets during reconstruction. The different colors represent different types of particles (ex: photons, charged hadrons, etc.)

Figure 3. Distribution of the number of primary vertices (per event) in data and MC in 2012.

There are three major classifications of pileup based upon the time at which the additional energy enters the calorimeter system. In-time (IT) pileup refers to energy from $pp$ collisions in the current bunch-crossing (BX) other than that at the hard scatter PV. This is the largest source of pileup energy. In addition, there is early out-of-time (EOOT) pileup, which refers to energy left in the calorimeters from previous BXs, and late out-of-time (LOOT) pileup, which refers to energy from later BXs that is integrated with the current event’s energy.

2. Pileup Measurement
While pileup itself cannot be directly measured, it can be correlated to various other directly measurable quantities. As pileup comes from additional $pp$ interactions, the number of PV ($N_{PV}$) is directly correlated to the amount of pileup; the greater the $N_{PV}$, the more pileup energy is added to the jets. The total offset for the event is equal to the average energy from a non-hard-scatter PV times the $N_{PV}$. The ability to correlate pileup with $N_{PV}$ is predicated on the CMS detector having a high reconstruction efficiency for PVs. In creating the corrections we use the cuts $N_{dof} > 4$ and $IsFake = false$ to make sure we are measuring only good PV. Figure 3 shows the good agreement between the $N_{PV}$ in data and in MC [4].

Another quantity which is directly correlated with the amount of pileup energy is the median energy per area ($\rho$), which is calculated using the Fastjet algorithm [5]. In order to make this correlation, pileup must be approximated as a homogeneous noise which can be subtracted off.

3. Pileup Corrections
The JEC corrections at CMS can be separated into two categories, those done at the hardware level and those done in software. Detector level corrections can be timing changes, threshold level modifications, or any other change to the hardware that might reduce the actual amount of pileup energy within the jets during reconstruction. Software level corrections are performed after reconstruction and scale jet $p_T$ up or down based on the measured amount of pileup.

3.1. Detector Level
Detector level mitigation techniques have been extremely effective in reducing the amount of out-of-time (OOT) pileup. The barrel and endcap hadronic calorimeter system (HBHE) used four time-slices (TS=25 ns) as its integration time during 2011. The measured LOOT pileup contributed about 20% of the pileup energy with EOOT pileup contributing another 2-4%. In
2012, the integration time in the HBHE system was changed to two TS, which reduced the LOOT pileup to almost zero and also reduced the EOOT pileup. This change is depicted in figure 4. With four TS starting from TS 0, the energy that is integrated will contain all of the energy from the current event as well as a significant portion of the energy from the next event. The two TS scheme includes little to none of the LOOT pileup. This is a significant improvement as it not only reduces one of the contributors to the pileup to nearly zero, but it also means that the rest of the pileup energy is more easily correlated with the IT $N_{PV}$.

3.2. Software Level

Software level pileup mitigation techniques are essential to most analyses. The purpose of these offset correction techniques is to measure and remove the extra energy inside of the jets that is not associated with the high-$p_T$ vertex and its UE [1–4; 6].

3.2.1. Average Offset  

The average offset (AO) method seeks to determine the average amount of energy added to an event due to low-$p_T$ scatters (pileup). The correction is calculated using three different samples; a Zero-Bias data dataset, a Zero-Bias MC sample that is made by overlaying a neutrino gun sample with Minimum Bias events, and a QCD MC sample. In each of the three samples, the offset ($p_{T,offset}$) is measured on an event-by-event basis by finding the $p_T$ deposited in a jet with a certain cone size (usually 0.5 or 0.7 at CMS), centered at a specific $\eta - \phi$ location. The $\phi$ direction of the jet is chosen at random and a scan of $\eta$, in steps of 0.1 within $|\eta| < 5.0$, is performed for each event. For calorimeter (Calo) jets, the energy in the jet is the sum of the energy inside the calorimeter towers whose $\eta - \phi$ location is inside the radius of the jet cone. For particle flow (PF) jets, the energy is that of the sum of the PF candidates inside the jet cone [6]. This procedure is called the Random Cone Algorithm.

These results are then binned in 1-D histograms, one for each value in $|\eta|$, with $N_{PV}$ on the x-axis and the $<p_{T,offset}>$ on the y-axis. Figure 5 shows a sample plot of the AO versus $N_{PV}$ for PF jets in $2.0 < |\eta| < 2.1$. The offset, shown in figure 5, is modeled by:

$$ Offset(N_{PV}) = [0] + [1](N_{PV}) + [2](N_{PV})^2 $$  \hspace{1cm} (1)

As expected, the value at the y-intercept of the line ($N_{PV} = 0$) is the $p_T$ contribution of the OOT pileup. Whereas the first and second order terms in $N_{PV}$ are dominated by IT pileup, with a small contribution from OOT pileup.

The correction to the raw, uncorrected jet $p_T (p_T^{RAW})$ is then derived to be a multiplicative scale factor which, when applied to the $p_T^{RAW}$, removes the added pileup energy. The scale factor for the AO method is:

$$ C_{AverageOffset}(N_{PV}, p_T^{RAW}, \eta) = 1 - \frac{< Offset(N_{PV}, \eta)>}{p_T^{RAW}} $$  \hspace{1cm} (2)
Once applied, the scale factor should correct the $p_T^{RAW}$ of the jet back to the true (generator) $p_T$:

$$p_T^{COR} = p_T^{GEN} = C_{AverageOffset}(N_{PV}, p_T^{RAW}, \eta) \cdot p_T^{RAW}$$  \hspace{1cm} (3)

### 3.2.2. Hybrid Jet Area

When using the AO method, one assumes that every jet contains the same amount of additional energy from pileup. Therefore, the best that the AO corrections can do is to correction on an event-by-event basis. The jet area (JA) method takes into account variations in the energy from pileup on an event-by-event basis and on a jet-by-jet basis by using a new parametrization [5; 7]. The hope is to improve the jet energy resolution and gain adaptability for future pileup conditions.

#### Figure 5. The average PF jet offset as a function of $N_{PV}$ within $2.0 < |\eta| < 2.1$.

#### Figure 6. Distribution of $p_T/A$ for all of the jets in a single event.

#### Figure 7. Distribution of $\rho$ versus $N_{PV}$ in data.

The new parametrization was developed based on the Fastjet algorithm. First the area ($A_j$) is calculated for each jet in the event. The next quantity needed is the median energy density of an event, $\rho$. This quantity is defined as the median of the distribution of $p_{Tj}/A_j$, where $j$ is the jet index. The $\rho$ computation is uses the $k_T$ jet clustering algorithm (R=0.6) because of its well-defined energy clustering and is insensitive to jets from a hard scatter as seen in figure 6. The sparsely populated values around $p_T/A = 30$ are the hard jets in the event (higher $p_T$ per unit area), whereas the large clump of jets lower in the spectrum represent the jets from pileup.

While the JA method is more granular and can correct on a jet-by-jet basis, the Fastjet algorithm used to parametrize such corrections does not know that the response of the detector is $\eta$ dependent. In the AO corrections, this was taken into account by binning the corrections in $\eta$. Therefore, a method was developed that combines the $\eta$ parametrization benefit of the AO method with the jet-by-jet granularity and parametrization of the JA method. This new method was called the hybrid jet area (HJA) method.

We start by parametrizing $\rho$ in terms of $N_{PV}$, as is done in figure 7. Once this relationship is obtained, the formula is inverted to obtain $N_{PV}(\rho)$. This relationship is then inserted into the function for the offset derived above.

$$Offset(N_{PV}(\rho)) = [0] + [1](N_{PV}(\rho)) + [2](N_{PV}(\rho))^2$$  \hspace{1cm} (4)

This function has all of the same characteristics, irrespective of the parametrization, as the
equation derived by the AO method. However, the HJA scale factor

\[ C_{Fastjet-based}(\rho, A, p_T^{\text{RAW}}, \eta) = 1 - A \frac{\text{Offset}(\rho, \eta)}{p_T^{\text{RAW}}} \]  

uses the JA to correct on a jet-by-jet basis. As seen in figure 8, both of these methods measure the pileup energy almost equivalently, on average [4].

![Figure 8. Comparison of the pileup measured in PF jets using AO method and the HJA method.](image)

![Figure 9. The particle flow composition of an \( N_{PV} = 1 \) event in MC.](image)

### 3.2.3. Charged Hadron Subtraction

Charged hadron subtraction (CHS) is an additional pileup subtraction tool that can be used in conjunction with either of the methods above, but which requires the use of particle flow jets. CHS uses the CMS detector’s excellent tracking capabilities to identify and remove jet constituents (charged hadrons) which are known to have originated from pileup vertices. The algorithm does not remove unassociated tracks, as they may have originated from the high-\( p_T \) vertex.

The benefit to applying this additional algorithm to the jet corrections is that it lowers the amount of pileup energy that the AO or HJA methods must compensate for, thus lowering the multiplicative scale factors. Furthermore, CHS has the additional advantage of being a particle-by-particle pileup subtraction technique, while at best the other algorithms work on an event-by-event or jet-by-jet basis.

### 3.3. Closure

In order to determine the efficacy of the corrections described above, we study the response of the detector to the jets, defined as \( R = \langle p_T^{\text{COR}}/p_T^{\text{GEN}} \rangle \). If the pileup corrections were ideal, we would expect the response for the jets to be identical at all \( N_{PV} \), meaning that there was no response dependence upon the pileup conditions inside the detector. In figure 10 the left plot shows the response as a function of the GenJet \( p_T \) for PF jets. The plot clearly shows that the response for low \( p_T \) jets is not equal to 1 (R=1 being a perfect response). In addition, the
response varies significantly for differing values of $N_{PV}$. However, the right plot shows the same thing after the pileup corrections have been applied. One can see not only a better response at low $p_T$, but that there is no $N_{PV}$ (pileup) dependence. Because all of the pileup energy has been removed from the jets, they also have a significantly lower response at low $p_T$. Any non-closure still present is due to various other effects, such as the detector’s $\eta$ and $p_T$ dependencies, which are corrected for using additional steps in the factorized jet correction chain [3; 4].

4. Conclusion

All of the various pileup removal techniques discussed in this paper have one goal, remove the dependence on pileup energy from the jet response. The first level of mitigation against pileup comes during the detector level readout and reconstruction efforts (ex: time-slices in the HBHE system). These have been shown to provide significant improvements in reducing OOT pileup.

The second level of pileup mitigation was at the software level and came in several of varieties. The AO method is the recommended correction for Calo jets as there is empirical evidence for slightly higher resolution in Calo jets when using this method. The Fastjet-based, HJA method is used for most analyses in CMS and is the recommended correction for PF jets. The key abilities of this method were its adaptability and per event/jet granularity. CHS is an additional tool which is very promising when used in conjunction with the other pileup mitigation techniques.

Once applied, there is good closure over most of the jet $p_T$ spectrum (in MC). Any remaining non-closure in the jet response is well understood and is removed at later stages of the factorized JEC chain [6].

References

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