Ultra-low Doping on Two-Dimensional Transition Metal Dichalcogenides using DNA Nanostructure Doped by a Combination of Lanthanide and Metal Ions

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Here, we propose a novel DNA-based doping method on MoS2 and WSe2 films, which enables ultra-low n- and p-doping control and allows for proper adjustments in device performance. This is achieved by selecting and/or combining different types of divalent metal and trivalent lanthanide (Ln) ions on DNA nanostructures, using the newly proposed concept of Co-DNA (DNA functionalized by both divalent metal and trivalent Ln ions). The available n-doping range on the MoS2 by Ln-DNA is between $6 \times 10^9$ and $2.6 \times 10^{10}$ cm$^{-2}$. The p-doping change on WSe2 by Ln-DNA is adjusted between $-1.0 \times 10^{10}$ and $-2.4 \times 10^{10}$ cm$^{-2}$. In Eu$^{3+}$ or Gd$^{3+}$-Co-DNA doping, a light p-doping is observed on MoS2 and WSe2 ($-10^{10}$ cm$^{-2}$). However, in the devices doped by Tb$^{3+}$ or Er$^{3+}$-Co-DNA, a light n-doping ($-10^{10}$ cm$^{-2}$) occurs. A significant increase in on-current is also observed on the MoS2 and WSe2 devices, which are, respectively, doped by Tb$^{3+}$- and Gd$^{3+}$-Co-DNA, due to the reduction of effective barrier heights by the doping. In terms of optoelectronic device performance, the Tb$^{3+}$ or Er$^{3+}$-Co-DNA (n-doping) and the Eu$^{3+}$ or Gd$^{3+}$-Co-DNA (p-doping) improve the MoS2 and WSe2 photodetectors, respectively. We also show an excellent absorbing property by Tb$^{3+}$ ions on the TMD photodetectors.

Deoxyribonucleic acid (DNA) nanostructures are currently being considered one of the promising materials for next-generation nanotechnology owing to the self-assembly and highly selective binding properties of DNA, which are based on the Watson-Crick complementary rules1–5. The self-assembled DNA nanostructures have been recently used in various nanoscale research fields, such as spintronics6–7, nanoelectronics8,9, biosensors10,11, and nanophotonics12,13. In particular, due to the highly selective binding property of DNA nanostructures, various particles such as metal nanoparticles (NPs)14,15, protein molecules16, nanowires (NWs)17, and divalent metal ions18–21 have been successfully incorporated onto the bases and backbone sites of DNA, eventually functionalizing the DNA nanostructures. Braun et al.22 and Maune et al.23 also reported self-assembled single-walled carbon nanotubes (SWNTs) bound on natural and origami DNA templates.

 Meanwhile, transition metal dichalcogenides (TMDs) with a two-dimensional layered structure, such as molybdenum disulfide (MoS2) and tungsten diselenide (WSe2), have been considered highly promising materials for next-generation flexible, wearable, stretchable and transparent devices due to their unique physical, electrical and optical properties24–30. TMD-based electronic devices are immune to short-channel effects owing to their

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centrations of Ln ions were added into the DX lattice solution (0, 0.5, 1 and 2 mM of Eu

suitable doping technique because precise control of the threshold voltage (V_{TH}) and the number of tightly-bound trions as well as very low contact resistance are required to achieve high performance. It is especially critical to develop an ultra-low level doping technique for the proper design and optimization of TMD-based devices (V_{TH} for transistors and tightly-bound trions for photodetectors) because high level doping (about 10^{12} cm^{-2}) causes TMD to act as a near-metallic layer. However, it is difficult to apply an ion implantation technique to TMD materials to achieve low-level doping due to crystal damage that occurs during the implantation process. Although safe doping techniques have recently been developed which can be applied to TMDs with an atomic layer structure, most of the previous TMD doping techniques employing potassium, Au NPs, polyethyleneimine (PEI), functional self-assembled monolayers (SAMs), and plasma treatment with fluorine (F) or oxygen (O) presented very high doping levels of 10^{12} cm^{-2}. Recently, low-level n- and p-doping of TMD materials was achieved using cesium carbonate (Cs_{2}CO_{3}), octadecyltrichlorosilane (OTS) and M-DNA, but further studies are needed to reduce the doping level down to an intrinsic level.

Here, we demonstrate a novel DNA-based doping method on selected 2D TMD materials (MoS_{2} and WSe_{2}), which represent n- and p-channel materials, respectively, enabling ultra-low-level n- and p-doping control, which had not been previously accomplished. This is achieved by selecting and/or combining different types of positive ions on DNA nanostructures. The DNA is functionalized by trivalent lanthanide ions (Ln-DNA) or both divalent metal and trivalent lanthanide ions (Co-DNA), which is newly proposed in this work. Since the phosphate backbone (PO_{4}^{3-}) in DNA causes n-doping on the TMD films by attracting and holding hole carriers on the side of TMD, various selections and/or combinations of divalent and trivalent positive ions on DNA make it possible to achieve both very low-level n- and p-doping, which are very close to the intrinsic level. First of all, Ln-DNA nanostructures based on double-crossover (DX) DNA lattices are explored as a function of ion concentration (here, Gd^{3+}) through Raman spectroscopy and atomic force microscopy (AFM) analyses. We then investigate low-level n- and p-doping phenomena on MoS_{2} and WSe_{2} by Ln-DNA and Co-DNA in terms of the performance (threshold voltage, on-/off-currents, photoresponsivity, and detectivity) of electronic and optoelectronic devices using Raman spectroscopy, photoluminescence (PL) spectroscopy, and electrical/optical measurements (I_{P} – V_{G} with/without exposure to 655-nm, 785-nm, and 850-nm lasers). We also discuss the absorption of trivalent lanthanide (Ln) ions with/without DNA templates with respect to the specific wavelengths of light and its influence on the performance of Ln-DNA or Co-DNA-doped TMD optoelectronic devices.

Results and Discussion

Synthesis and analysis of Ln-DNA nanostructures. As shown in Fig. 1a, we formed a DNA double-crossover (DX) lattice structure with various concentrations of lanthanide ions and then investigated its stability by Raman spectroscopy and AFM to avoid structural deformation of the DX lattices. The DNA DX lattices were fabricated by a conventional solution-free annealing process. After annealing at 95°C, various concentrations of Ln ions were added into the DX lattice solution (0, 0.5, 1 and 2 mM of Eu^{3+}, Gd^{3+}, Tb^{3+}, and Er^{3+}). In order to verify the stability of Gd-DNA, which is Ln-DNA with various concentrations of Gd^{3+}, we obtained Raman spectra of Gd-DNA grown on glass substrate (Fig. 1b) and identified peak differences between 0 mM (pristine DNA), 0.5 mM, 1 mM, and 2 mM of [Gd^{3+}]. In the pristine DNA crystal without Gd^{3+}, Raman peak signals were observed at 1246 and 1420 cm^{-1} for adenine (A), 770 and 1469 cm^{-1} for thymine (T), 931 and 1590 cm^{-1} for guanine (G), 618, 655, and 1348 cm^{-1} for cytosine (C), and 1066 and 1145 cm^{-1} for the phosphate backbone (PO_{4}^{3-})40,41. After binding Gd^{3+} ions onto the base pairing and backbone sites, the Raman intensity related to base and backbone sites was reduced. In particular, a Raman peak shift was observed in the base pairing sites only when the concentration of [Gd^{3+}] was above 0.5 mM. Because the Ln ions (including Gd^{3+}) were reported to be normally attached to backbone sites due to its intrinsic radius (300 ~ 340 pm), which is larger than the hydrogen bonding length (340 pm) in DNA bases, it is thought that the Ln ions are intercalated in the bases without chemical bonding, and therefore seem to influence the DNA nanostructure. As the concentration of [Gd^{3+}] increased, the Raman peak intensity of Gd-DNA decreased, and it was eventually hard to identify Raman peaks of Gd-DNA above 1 mM of [Gd^{3+}] indicating that the Gd-DNA crystal structure is deformed when more than 1 mM of [Gd^{3+}] ions are added. This was verified once again through AFM images of Gd-DNA with 0 mM (pristine DNA), 0.5 mM, and 2 mM of [Gd^{3+}] in Fig. 1c. The AFM images of pristine and Gd (0.5 mM) samples showed polycrystalline structures without deformation, and the clear periodicity of each DX lattice was also confirmed in the noise-filtered 2D spectrum image after fast Fourier transform (FFT). However, in the case of 2 mM of [Gd^{3+}], Gd-DNA complexes seemed to be aggregated, and consequently form an amorphous structure. A similar phenomenon was also observed in the cases of other Ln ions with different concentrations, and the determined optimum concentrations of each Ln ion were 1 mM for [Eu^{3+}], 1 mM for [Gd^{3+}], 1 mM for [Tb^{3+}], and 1 mM for [Er^{3+}].

Raman analysis of TMD films doped by Ln-DNA or Co-DNA. The artificially designed Ln- or Co-DNA DX solution was dropped and dried five times on the MoS_{2} and WSe_{2} surfaces. DNA nanostructures were homogeneously dispersed along the surface of MoS_{2} and WSe_{2} because of the self-aligning capability of the DNA39. In the previous study, the PO_{4}^{3-} backbone sites in DNA nanostructures were reported to induce and hold positive charges (holes) at the interface region in the side of MoS_{2} and WSe_{2}; because of the self-aligning capability of the DNA39. In the previous study, the PO_{4}^{3-} backbone sites in DNA nanostructures were reported to induce and hold positive charges (holes) at the interface region in the side of MoS_{2} and WSe_{2} because of the self-aligning capability of the DNA39. In the previous study, the PO_{4}^{3-} backbone sites in DNA nanostructures were reported to induce and hold positive charges (holes) at the interface region in the side of MoS_{2} and WSe_{2} because of the self-aligning capability of the DNA39. In the previous study, the PO_{4}^{3-} backbone sites in DNA nanostructures were reported to induce and hold positive charges (holes) at the interface region in the side of MoS_{2} and WSe_{2} because of the self-aligning capability of the DNA39.

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analysis was then performed on the Ln-DNA-doped TMD samples to investigate the Ln-DNA doping effects by Eu\(^{3+}\), Gd\(^{3+}\), Tb\(^{3+}\), and Er\(^{3+}\) ions on TMD films, as shown in Supplementary Information Fig. S3. In the MoS\(_2\) films, two conventional peaks (E\(_{1g}^{1}\) and A\(_{1g}\)) were observed at ~380 cm\(^{-1}\) and ~406 cm\(^{-1}\), which respectively indicate the in-plane and out-of-plane vibrations for bulk MoS\(_2\)\(^{43}\). On the other hand, only a single peak was obtained at ~250 cm\(^{-1}\) in WSe\(_2\) films because both the E\(_{1g}^{1}\) and the A\(_{1g}\) modes for WSe\(_2\) are close to 250 cm\(^{-1}\)\(^{44}\). To clarify the degree of Ln-DNA doping by each ion on TMD films, we extracted the peak shift values in E\(_{1g}^{1}\), A\(_{1g}\), and E\(_{1g}^{1}\) + A\(_{1g}\) peaks which have respectively peak shift values in the ranges of ~1.3 ~ ~0.8 cm\(^{-1}\), ~1.2 ~ ~0.7 cm\(^{-1}\), and ~1.3 ~ ~0.9 cm\(^{-1}\), indicating that very low level n-type doping was achieved on the TMD films. This n-type doping phenomenon by Ln-DNA also seems to be slightly weaker than that by pristine DNA. For reference, a blue-shift phenomenon in all of the peaks was observed after the M-DNA doping on the MoS\(_2\) and WSe\(_2\) films (2.8 ~ 5.0 cm\(^{-1}\) for E\(_{1g}^{1}\), 3.8 ~ 5.2 cm\(^{-1}\) for A\(_{1g}\) and 3.1 ~ 5.1 cm\(^{-1}\) for E\(_{1g}^{1}\) + A\(_{1g}\))\(^{39}\), indicative of p-type doping.

Next, we added Cu\(^{2+}\) ions with a relatively smaller ionic radius into Ln-DNA to acquire weaker n- or p-type doping effects on TMD films by converting the negative charges at the base pairings to positive charges; these
new DNA nanostructures are referred to as Co-DNA. Here, the Cu$^{2+}$ ions are expected to be bound at DNA base pairings because the Ln ions are predominantly attached to the backbone sites, eventually modulating the strength of the total charge which previously showed weak negative polarity. In this Co-DNA experiment, the different amounts of Ln ions (0.5 mM Eu$^{3+}$, 0.5 mM Gd$^{3+}$, 0.5 mM Tb$^{3+}$, or 0.5 mM Er$^{3+}$) were mixed with 2 mM of Cu$^{2+}$ ions to avoid structural deformation of Co-DNA nanostructures. Then, Raman spectroscopy was performed on the DNA (pristine), Ln-DNA (by Gd$^{3+}$), and Co-DNA (by Gd$^{3+}$ and Cu$^{2+}$) structures, as shown in Fig. 3b. As previously mentioned, the intensity of Raman peaks, indicating the base pairing and backbone sites of DNA, was reduced in the case of Gd-DNA due to the Gd$^{3+}$ ions attached on the sites. When Cu$^{2+}$ ions were added into the Gd-DNA, we observed an additional reduction in all of the peaks related to base pairings and backbone sites, indicating that Cu$^{2+}$ ions were additionally attached on the remaining backbone sites as well as the base pairings; it was previously predicted that Ln ions were not bound to the base pairings. Finally, we coated the various Co-DNA nanostructures on MoS$_2$ and WSe$_2$ films, and performed Raman analysis on the films. The raw Raman peak data of Co-DNA-doped MoS$_2$ and WSe$_2$ can be found in Supplementary Information Fig. S3. Figure 3c,d show the extracted peak shift values before/after Co-DNA doping for MoS$_2$ ($\Delta E_{2g}$ and $\Delta A_{1g}$) and WSe$_2$ ($\Delta E_{2g}$,$\Delta A_{1g}$), respectively. Here, the Raman measurement was performed on fifty different points in each sample. For both the MoS$_2$ and WSe$_2$ cases, Co-DNA nanostructures with Tb$^{3+}$ or Er$^{3+}$ ions seem to induce much weaker n-type doping phenomena compared to the previous Tb- and Er-DNA (Ln-DNA). This is because smaller peak shift values were observed in the peaks ($-0.5 \sim -0.7$ cm$^{-1}$ for $E_{2g}$, $-0.3 \sim -0.5$ cm$^{-1}$ for $A_{1g}$ and $-0.2 \sim -0.4$ cm$^{-1}$ for $E_{2g}$+$A_{1g}$). Meanwhile, a blue-shift phenomenon was confirmed in the MoS$_2$ and WSe$_2$ films doped by Co-DNA with Eu$^{3+}$ or Gd$^{3+}$ ions, indicating that the films were lightly p-doped. Here, the peak shift values are $0.3 \sim 0.1$ cm$^{-1}$ for $E_{2g}$, $0.6 \sim 0.4$ cm$^{-1}$ for $A_{1g}$, and $0.45 \sim 0.4$ cm$^{-1}$ for $E_{2g}$+$A_{1g}$, which are much smaller than the previously reported values of M-DNA-doped films$^{24}$. These results support the claim that the Cu$^{2+}$ ions are bound at DNA base pairings (also, remained backbone sites) and can be used to modulate the strength of total charge in Ln-DNA, eventually making the Co-DNA nanostructure show weak negative or positive polarity. In particular, based on the ionic radius and optimum concentrations shown in Table 1, it is thought that Eu$^{3+}$ ions show the strongest positive charge due to their largest ionic radius and highest concentration. As expected, the largest peak shift was observed in the MoS$_2$ and WSe$_2$ films doped by Co-DNA with Eu$^{3+}$ ions. In contrast to the Eu$^{3+}$ sample, the smallest peak shift, indicating light n-type doping, was obtained in the TMD films doped by Co-DNA with Er$^{3+}$ ions which are expected to show the lowest positive charge strength owing to the smallest ionic radius.
Electrical characteristics of Ln-DNA- and Co-DNA-doped TMD electronic devices. We then fabricated TMD electronic devices (back-gate transistors) and investigated the electrical properties (threshold voltage, 2D sheet carrier concentration, and on-current) of TMD devices before/after Ln-DNA or Co-DNA doping.

Figure 3. Doping mechanism by Co-DNA and Raman analysis of TMD films doped by Co-DNA. (a) Schematic diagrams explaining Co-DNA doping on TMD. (b) Raman spectra of DNA, Gd³⁺-based Ln-DNA (or Gd-DNA), and Gd³⁺-based Co-DNA (with Cu²⁺ ions). (c) Extracted Raman peak shift data of MoS₂ and WSe₂ doped by Co-DNA.

|               | Ln-DNA | Co-DNA |
|---------------|--------|--------|
| DNA           |        |        |
| Gd³⁺-based    |        |        |
| Gd-DNA        |        |        |
| Gd³⁺-based Co-DNA (with Cu²⁺ ions) |        |        |
| DNA           | T      | T      |
| Gd³⁺-based    | G      | G      |
| Gd-DNA        | T      | T      |
| Gd³⁺-based Co-DNA (with Cu²⁺ ions) | T      | T      |

Table 1. Atomic numbers, ionic radii, and optimum concentrations of various Ln ions.

|               | Eu | Gd | Tb | Er |
|---------------|----|----|----|----|
| Atomic number | 63 | 64 | 65 | 68 |
| Ionic radius (pm) | 108.7 | 107.8 | 106.3 | 103 |
| Optimum concentration (mM) | 1 | 1 | 1 | 1 |

We then fabricated TMD electronic devices (back-gate transistors) and investigated the electrical properties (threshold voltage, 2D sheet carrier concentration, and on-current) of TMD devices before/after Ln-DNA or Co-DNA doping.
Figure 4 shows the schematic diagram of TMD transistors and the energy band diagrams of metal-undoped/doped TMD junctions in the source-side when proper operating biases (positive $V_{DS}$ for MoS$_2$ and negative $V_{DS}$ for WSe$_2$) are applied. When Ln-DNA is coated on TMD films, as already explained, the negative charges of Ln-DNA are expected to hold hole carriers at the interface between Ln-DNA and TMD, resulting in an n-type doping phenomenon. In the case of MoS$_2$, as shown in Fig. 4b, the electric field at the source-side Ti-MoS$_2$ junction is predicted to increase due to the down-shift of MoS$_2$ energy bands by Ln-DNA doping, thereby reducing the effective barrier height of the Ti-MoS$_2$ junction (Schottky barrier lowering effect: $\Phi_{Control,eff} > \Phi_{Ln^{3+},eff}$). As a result, more electron carriers may be injected from Ti to MoS$_2$, and this result in a negative shift in $V_{TH}$. In addition, it is also predicted that the WSe$_2$ energy band shifts down after Ln-DNA doping. Subsequently, the effective hole barrier height from the source to WSe$_2$ increases, finally causing a negative shift in $V_{TH}$. On the other hand, because Co-DNA shows very weak positive polarity (Eu$^{3+}$ or Gd$^{3+}$) or negative polarity (Tb$^{3+}$ or Er$^{3+}$), it is possible to hold electron or hole carriers at the interface, respectively, eventually causing very low level n- or p-type doping. In particular, the level of n-doping by Tb- or Er-based Co-DNA is expected to be lower than that by Ln-DNA. In MoS$_2$, Eu- or Gd-based Co-DNA moves up the MoS$_2$ energy band, consequently reducing the electric field at the source-side junction and increasing the effective electron barrier height (green line; a positive shift in $V_{TH}$). In contrast, Tb- or Er-based Co-DNA causes n-doping on MoS$_2$, thereby shifting down its energy band and eventually decreasing the effective electron barrier height (pink line; a negative shift in $V_{TH}$). A similar phenomenon in the effective hole barrier height and $V_{TH}$ is also expected on WSe$_2$, as shown in the right-bottom of Fig. 4b.

Figure 4. Schematic diagram of back-gated TMD doped device and Ln- or Co-DNA doping effect on TMD films. (a) Schematic diagram showing the back-gated transistor (or photodetector) device fabricated on TMD films doped by Ln- or Co-DNA. (b) Energy band diagrams of source-undoped/doped TMD junctions.
Figure 5. Electrical characterization of TMD transistors doped by Ln-DNA. $I_D - V_G$ characteristics of the (a) MoS$_2$ and (d) WSe$_2$ transistors undoped/doped by Gd$^{3+}$-DNA. Threshold voltage shifts ($\Delta V_{TH} = V_{TH,Ln-DNA} - V_{TH,Control}$) and variations of carrier concentration ($\Delta n = n_{Ln-DNA} - n_{Control}$ and $\Delta p = p_{Ln-DNA} - p_{Control}$) extracted in (b) MoS$_2$ and (e) WSe$_2$ transistors, which were undoped/doped by Ln-DNA. On-current ratio ($I_{on, Ln-DNA}/I_{on, Control}$) extracted in the undoped/doped (c) MoS$_2$ and (f) WSe$_2$ transistors.

Figure 5a shows $I_D - V_G$ characteristics of undoped and Gd-DNA-doped MoS$_2$, which were measured at $V_{DS} = 5$ V. As previously predicted, a negative shift in $V_{TH}$ (from $-28.5$ V to $-31.2$ V) and an increase in on-current (from $2.3 \times 10^{-6}$ A/µm to $1.2 \times 10^{-5}$ A/µm) were observed after the Gd-DNA doping, indicating that the Gd-DNA caused n-type doping of MoS$_2$. These changes are attributed to an increase in the tunneling probability from the source to the MoS$_2$ channel by the Gd-DNA doping. These changes in $V_{TH}$ and on-current were also observed in the other Ln-DNA with Eu$^{3+}$, Tb$^{3+}$, or Er$^{3+}$ ions, as shown in Fig. 5b,c. The extracted $\Delta V_{TH} (= V_{TH,Ln-DNA} - V_{TH,Control})$ values were between $-1$ and $-2.5$ V, indicating the n-doping of MoS$_2$, which is also consistent with the previous Raman analysis. The difference in 2D sheet doping concentrations ($\Delta n = n_{Ln-DNA} - n_{Control}$) extracted in the Ln-DNA-doped MoS$_2$ devices also showed positive values, which mean an increase in the number of electron carriers and consequently in n-doping of MoS$_2$. The lowest $\Delta n$ was $-6 \times 10^{-6}$ cm$^{-2}$ in the case of Tb-DNA, and the available n-doping range by Ln-DNA was between $-6 \times 10^{-6}$ cm$^{-2}$ to $-2.6 \times 10^{-5}$ cm$^{-2}$, in which the device performance can be controlled. In addition, we confirmed the improvements in on-current through Ln-DNA doping in Fig. 5c. The on-current ratio ($= I_{on,Ln-DNA}/I_{on,Control}$) was between approximately 8 and 10. This enhancement is thought to be caused by the reduction in effective electron barrier height through Ln-DNA doping, indicating that there was a reduction in contact resistance. Based on simple on-state contact resistance extraction$^{45}$, Gd-DNA doping reduced the estimated contact resistance to $-72.2$ kΩ (from $-90.7$ kΩ) on the MoS$_2$ device and increased it to $95.1$ kΩ (from $-67.7$ kΩ) on the WSe$_2$ device. On the other hand, we observed a degradation in the electronic performance of WSe$_2$ devices after Ln-DNA doping, as shown in Fig. 5d, which presents the $I_D - V_G$ characteristics of undoped/doped WSe$_2$ devices by Gd-DNA (at $V_{TH} = -5$ V). Here, a negative shift in $V_{TH}$ (from 10 V to 8 V) and a decrease in on-current (from $5.2 \times 10^{-5}$ A/µm to $9.2 \times 10^{-5}$ A/µm) were confirmed, indicating the n-doping of WSe$_2$. Then, $\Delta V_{TH}$ values extracted in each Ln-DNA-doped WSe$_2$ device were between $-1.4$ V and $-3.1$ V; we also obtained negative $\Delta p$ values ($= p_{Ln-DNA} - p_{Control}$) in the range from $-1.0 \times 10^{10}$ cm$^{-2}$ to $-2.4 \times 10^{10}$ cm$^{-2}$ (Fig. 5e). These changes are consistent with the red-shift phenomenon in the Raman peaks mentioned above, which is indicative of n-doping (reduction of hole concentration in WSe$_2$). In contrast to the improvement in on-current in MoS$_2$ devices, Ln-DNA doping degraded the parameters of WSe$_2$ devices, as shown in Fig. 5f. The on-current ratio was found to be below 0.3, likely because the energy band of WSe$_2$ was shifted down after Ln-DNA doping. Consequently, its effective hole barrier height from the source to WSe$_2$ was increased, resulting in an increase in contact resistance.

However, as already expected, slightly different changes were observed on the TMD-based electronic devices after doping by the Co-DNA. Figure 6a shows the $I_D - V_G$ characteristics of MoS$_2$ devices undoped and doped by Co-DNA (Gd$^{3+}$ + Cu$^{2+}$), where a positive shift in $V_{TH}$ (from $-30$ V to $-28.2$ V) and a reduction in on-current (from $8.2 \times 10^{-5}$ A/µm to $4.2 \times 10^{-5}$ A/µm) were observed after the Gd$^{3+}$-based Co-DNA doping, indicating a
light p-doping of MoS$_2$. Based on the extracted $\Delta V_{TH}$ ($= V_{TH_{Co-DNA}} - V_{TH_{Control}}$) and variations of carrier concentration ($\Delta n = n_{Co-DNA} - n_{Control}$ and $\Delta p = p_{Co-DNA} - p_{Control}$) extracted in (b) MoS$_2$ and (e) WSe$_2$ transistors, which were undoped/doped by Co-DNA. On-current ratio ($I_{on\ ratio} = I_{on_{Co-DNA}}/I_{on_{Control}}$) extracted in the undoped/doped (c) MoS$_2$ and (f) WSe$_2$ transistors.

As seen in Fig. 6f, the on-current ratio of the WSe$_2$ devices before/after doping was also increased up to 4.1 when Eu$^{3+}$- and Gd$^{3+}$-based Co-DNA doping (p-doping) was performed. However, the on-current ratio was reduced below 0.25 in devices doped with Tb$^{3+}$- and Er$^{3+}$-based Co-DNA (n-doping). In additional experiments on tri-layer TMDs, we confirmed that the doping effects of Ln- or Co-DNA were independent of TMD...
and 0.17–0.23 V for the buffer solution). In addition, the low dark current level. We then extracted and plotted photoresponsivity (R) of the electron channel is improved as the thickness. We also note that the influence of four kinds of solvents used during the fabrication process (DI water, acetone, IPA, and buffer solutions) on TMD-based devices (here, ΔVTH) seems to be negligible when compared to the effect of Ln- or Co-DNA-doping (0.3–0.33 V for DI water, 0.2–0.22 V for acetone, −0.24—−0.2 V for IPA, and 0.17–0.23 V for the buffer solution). In addition, the ΔVTH values of the Ln- and Co-DNA-doped MoS2 and WSe2 devices changed slightly after 120 hours of air-exposure (Supplementary information Fig. S7). The ΔVTH values of Ln-DNA-doped MoS2 and WSe2 transistors increased 15–20%, indicating the weakening of the n-type doping effects after 120 hours. The ΔVTH values of the Co-DNA-doped samples also decreased (Eu- or Gd-based Co-DNA) or increased (Tb- or Er-based Co-DNA) as a function of air exposure time because of the reduced dipole moment of the phosphate backbone (PO4−) and lanthanide ions (Ln3+) related to the humidity-associated structural deformation of Ln- and Co-DNA nanostructures46.

Characteristics of Co-DNA-doped TMD optoelectronic devices. We then performed a photocurrent measurement on the Co-DNA-doped TMD photodetector devices with 655-nm, 785-nm, and 850-nm lasers in order to investigate the effects of p- and n-doping by Co-DNA on optoelectronic device performance. Figure 7a shows a schematic diagram of the Co-DNA-doped TMD photodetector, along with the corresponding energy band diagrams of i) MoS2 and ii) WSe2 before/after Co-DNA doping. Compared to the control device, the depletion width at the source/MoS2 junction is expected to be broadened after Eu3+ or Gd3+-based Co-DNA doping (pink solid line: p-doping), consequently helping to collect photocarriers and increase photocurrent. However, the n-doping phenomenon by Tb3+ or Er3+-based Co-DNA is predicted to narrow the depletion width and reduce photocurrent below that of the control device (green dotted line: n-doping). In contrast, WSe2 photodetectors doped by the Tb3+ or Er3+-based Co-DNA show higher photocarrier collection and subsequently higher photocurrent. Figure 7b shows Iph–VGS characteristics of a Ga3+-based Co-DNA-doped TMD photodetector before and after light exposure. In the off-state (VGS = VTH), the absence of majority carrier paths in MoS2 (low electron current) and the high hole barrier height (low hole current) are expected to reduce the dark current below a level of 10−11A/μm. As a result, the photocurrent was more clearly observed in the off-state because of the low dark current level. We then extracted and plotted photoresponsivity (R = Iphoton/Idark) and detectivity (D* = (RA/(TΔVTH)) values as a function of VGS − VTH in Fig. 7c,d. Here, the generated photocurrent is Iphoto−Idark, T is the total incident optical power, A is the effective area of the device, and is the absolute value of the electron charge (1.6 × 10−19C), and Idark is dark current. In Fig. 7c, Eu3+ or Gd3+-based Co-DNA-doped MoS2 devices showed higher photoresponsivity (~68μA/W at VGS = VTH) than those (~25μA/W at VGS = VTH) of Tb3+ or Er3+-based Co-DNA samples because of their broader depletion width and subsequent higher photocurrent. The photoresponsivity was increased as a function of VGS in all cases of Co-DNA because the conductivity of the electron channel is improved as VGS approaches VTH. An opposite trend was observed in detectivity (D*) because D* is expressed as the ratio of Iphoto to Idark and is also affected by the Idark, as shown in Fig. 7d. The highest value (maximum D* = ~6.02 × 109 Jones at VGS = VTH = −2.6 V) was obtained in the Eu3+-based Co-DNA-doped device. Figure 7e shows Iph–VGS characteristics of the WSe2 device doped by Ga3+-based Co-DNA before/after laser exposure. Figure 7b shows photocurrents of (c) MoS2 and (f) WSe2 photodetectors as a function of VGS–VTH (in off-state). Detectivity of (d) MoS2 and (g) WSe2 photodetectors as a function of VGS–VTH (in off-state).

Figure 7. Schematic diagram and characterization of TMD photodetectors doped by various Co-DNA. (a) Schematic diagram of Co-DNA-doped MoS2/WSe2 photodetectors with a light source (λ = 785 nm and P = 0.6 mW), and the energy band diagrams of Ti–MoS2/Pt–WSe2 junctions under the illuminated condition. Iph–VGS characteristics of Ga3+-based Co-DNA-doped (b) MoS2 and (e) WSe2 photodetectors before/after the laser exposure. Photoresponsivity of (c) MoS2 and (f) WSe2 photodetectors as a function of VGS–VTH (in off-state).
light exposure, and we also extracted photoresponsivity and detectivity as a function of $V_{GS} - V_{TH}$ in Fig. 7f.g. An opposite trend was observed in photoresponsivity of WSe$_2$ devices compared to MoS$_2$ devices; that is, higher values were obtained in the n-doped devices by Tb$^{3+}$- and Er$^{3+}$-based Co-DNA. In the Er$^{3+}$-based Co-DNA-doped device, a maximum photoresponsivity of ~$27 \mu A/W$ was observed at $V_{GS} = V_{TH}$. In addition, like the MoS$_2$ photodetectors, the photoresponsivity of WSe$_2$ devices increased as the $V_{GS}$ moved to the $V_{TH}$ point. Detectivity was also better in the n-doped WSe$_2$ devices, and it increased as a function of $V_{GS} - V_{TH}$ where the maximum detectivity of $6.5 \times 10^4$ Jones was obtained in the Er$^{3+}$-based Co-DNA-doped device when $V_{GS} - V_{TH} = 3 \text{V}$ ($8.2 \times 10^4$ Jones at $V_{GS} - V_{TH} = 5.2 \text{V}$).

Then, the MoS$_2$ and WSe$_2$ photodetectors doped by Gd$^{3+}$-based Co-DNA were investigated under laser exposure with different wavelengths (655-nm, 785-nm, and 850-nm), as shown in Fig. 8. Figure 8a shows $I_D$ - $V_G$ characteristics of the TMD photodetectors measured before and after the various light exposures. In both devices an increase in photocurrent was observed as the laser wavelength decreased. The energy of incident light increases with decreasing wavelength, and it eventually overcomes the direct bandgap (at K valley) at $\lambda = 655$ nm (1.89 eV > 1.85 eV of direct $E_g$-MoS$_2$) and 1.65 eV of direct $E_g$-WSe$_2$), finally improving the absorption probability as seen in Fig. 8b. The photoresponsivity and detectivity ratio values ($R = R_{Co-DNA}/R_{Control}$ and $D^* = D_{Co-DNA}/D_{Control}$) were extracted from the $I_D$ - $V_G$ curves of the devices doped by Gd$^{3+}$ and Tb$^{3+}$-based Co-DNA, were then plotted as a function of wavelength in Fig. 8c,d, respectively. As already mentioned in Fig. 7, p-doping of MoS$_2$ by Gd$^{3+}$-based Co-DNA and n-doping of WSe$_2$ by Tb$^{3+}$-based Co-DNA presented better optoelectronic performance (higher photoresponsivity and detectivity) at 785-nm and 850-nm as compared to the control devices. After doping the MoS$_2$ and WSe$_2$ devices respectively with Tb$^{3+}$ and Gd$^{3+}$-based Co-DNA, worse performance was observed in the same wavelength range because of the narrowed depletion width. However, under exposure to a 655-nm laser, dramatic performance improvement was observed in the devices doped by Tb$^{3+}$-based Co-DNA, but not in the case of Gd$^{3+}$-based Co-DNA. This seems to be attributed to the excellent emission and absorption properties of Tb$^{3+}$-based Co-DNA around 655-nm, which were confirmed by PL analysis in Fig. 8e. Because of the Tb$^{3+}$ ions attached mainly on the backbone sites, much higher PL intensity was observed in the Co-DNA, where a broad PL spectrum consists of various PL peaks between 438-nm and 630-nm; their corresponding energy state levels are shown in Fig. 8f. Several research groups previously reported the superior emitting properties of Tb$^{3+}$ ions in the range from 380-nm to 700-nm. Through the energy state levels of Tb$^{3+}$ ions, the incident 655-nm light is expected to additionally generate electron-hole pairs in the Co-DNA region. After that, we predict that the photocarriers are transferred from the Co-DNA to the TMD layers, consequently increasing the number of total photocarriers and the photocurrent. We also note that Gd$^{3+}$ ions generally show emitting and absorbing properties in the wavelength range between 250-nm and 312-nm, which is far from the wavelengths of the lasers used.

**Conclusion**

In conclusion, we proposed a new DNA-based doping method on the MoS$_2$ and WSe$_2$ films (representative n- and p-channel materials, respectively), which enabled ultra-low n- and p-doping control by selecting and/or combining different types of positive ions on DNA nanostructures. Because the PO$_4$$^{3-}$ in DNA caused n-doping on the TMD films by attracting and holding hole carriers on the side of the TMD, the various selections and/or combinations of the divalent and trivalent positive ions on DNA made it possible to achieve both very low-level n- and p-doping, which are very close to the intrinsic level. The available n-doping range ($\Delta n$) on the MoS$_2$ by Ln-DNA was between $6 \times 10^9$ cm$^{-2}$ and $2.6 \times 10^{10}$ cm$^{-2}$, which was even lower than that by pristine DNA ($6.4 \times 10^9$ cm$^{-2}$). The p-doping change ($\Delta p$) on the WSe$_2$ by Ln-DNA was controlled from $1.0 \times 10^9$ cm$^{-2}$ to $2.4 \times 10^9$ cm$^{-2}$, which changes were consistent with the red-shift phenomenon in Raman peaks, which indicated n-doping (an increase of electron carriers in MoS$_2$ and a decrease of hole carriers in WSe$_2$). Then, the MoS$_2$ and WSe$_2$ photodetectors doped by Gd$^{3+}$-based Co-DNA, were then plotted as a function of wavelength in Fig. 8c,d, respectively. As already mentioned in Fig. 7, p-doping of MoS$_2$ by Gd$^{3+}$-based Co-DNA and n-doping of WSe$_2$ by Tb$^{3+}$-based Co-DNA presented better optoelectronic performance (higher photoresponsivity and detectivity) at 785-nm and 850-nm as compared to the control devices. After doping the MoS$_2$ and WSe$_2$ devices respectively with Tb$^{3+}$ and Gd$^{3+}$-based Co-DNA, worse performance was observed in the same wavelength range because of the narrowed depletion width. However, under exposure to a 655-nm laser, dramatic performance improvement was observed in the devices doped by Tb$^{3+}$-based Co-DNA, but not in the case of Gd$^{3+}$-based Co-DNA. This seems to be attributed to the excellent emission and absorption properties of Tb$^{3+}$-based Co-DNA around 655-nm, which were confirmed by PL analysis in Fig. 8e. Because of the Tb$^{3+}$ ions attached mainly on the backbone sites, much higher PL intensity was observed in the Co-DNA, where a broad PL spectrum consists of various PL peaks between 438-nm and 630-nm; their corresponding energy state levels are shown in Fig. 8f. Several research groups previously reported the superior emitting properties of Tb$^{3+}$ ions in the range from 380-nm to 700-nm. Through the energy state levels of Tb$^{3+}$ ions, the incident 655-nm light is expected to additionally generate electron-hole pairs in the Co-DNA region. After that, we predict that the photocarriers are transferred from the Co-DNA to the TMD layers, consequently increasing the number of total photocarriers and the photocurrent. We also note that Gd$^{3+}$ ions generally show emitting and absorbing properties in the wavelength range between 250-nm and 312-nm, which is far from the wavelengths of the lasers used.

**Experimental Methods**

**DNA DX lattice fabrication.** High-performance liquid chromatography (HPLC)-purified synthetic oligonucleotides of DNA was purchased from BIONEER (www.bioneer.com). Two DX tiles were used to construct a 2D DNA nanostructure through a conventional free solution annealing process. Complexes for the DX structure (200 nM) were formed by mixing a stoichiometric quantity of each strand in physiological $1 \times$ TAE/Mg$^{2+}$ buffer
Figure 8. Characterization of Co-DNA-doped TMD photodetectors under laser exposure with different wavelengths. (a) $I_D - V_G$ characteristics of MoS$_2$ and WSe$_2$ photodetectors doped by Gd$^{3+}$-based Co-DNA under laser exposure with different wavelengths (655, 785, and 850 nm). (b) Energy band structures of bulk MoS$_2$ and WSe$_2$. Here, we approximately marked the photon or phonon transition under the laser exposure with different wavelengths. (c) Photoresponsivity ratio ($R_{Co-DNA}/R_{Control}$) and (d) detectivity ratio ($D^*_{Co-DNA}/D^*_{Control}$) of Gd$^{3+}$ and Tb$^{3+}$-based Co-DNA-doped TMD photodetector under different laser exposures. (e) Photoluminescence data of DNA and Tb$^{3+}$-DNA crystals, and Gaussian fitting lines for Tb$^{3+}$-DNA PL data. Here, the excitation wavelength is 400 nm. (f) Energy state levels of Tb$^{3+}$-DNA and the schematic presenting the mechanism of photocarrier transfer from Tb$^{3+}$ ions to the TMD.
Pristine DNA and Tb3⁺, Er3⁺, or Cu2⁺. Here, the thickness of the gate oxide (SiO₂) was 90 nm.

**Electrical characterization of Ln-DNA or Co-DNA-doped TMD electronic/optoelectronic devices.** The fabricated transistor devices were electrically analyzed using an HP 4415B semiconductor parameter analyzer (Alpha300 M+, WITec). Here, TMD bulk flakes with similar thickness (~38 nm for MoS₂ and ~33 nm for WSe₂) were selected through AFM analysis in order to avoid the thickness effect. Raman spectroscopy with an excitation wavelength of 532 nm was used; the laser beam size was approximately 0.7 ~ 0.9 μm, and the instrumental spectral resolution was less than 0.9 cm⁻¹. An integration time of 5 seconds and a spectrometer with 1800 grooves/mm was employed for the test.

**Characterizations of Ln-DNA or Co-DNA-doped TMD films.** Ln-DNA or Co-DNA-doped TMD samples were investigated and compared with a control sample (undoped TMD) by PL/Raman spectroscopy (Alpha300 M+, WITec). Here, TMD bulk flakes with similar thickness (~38 nm for MoS₂ and ~33 nm for WSe₂) were selected through AFM analysis in order to avoid the thickness effect. Raman spectroscopy with an excitation wavelength of 532 nm was used; the laser beam size was approximately 0.7 ~ 0.9 μm, and the instrumental spectral resolution was less than 0.9 cm⁻¹. An integration time of 5 seconds and a spectrometer with 1800 grooves/mm was employed for the test.

**Fabrication of Ln-DNA or Co-DNA-doped TMD electronic/optoelectronic devices.** For the fabrication of back-gated TMD transistors, source/drain electrode regions were patterned (channel length and width are 5 μm) on MoS₂/SiO₂/Si samples by optical lithography, followed by 10-nm-thick Ti (for MoS₂) or Pt (for WSe₂) and 50-nm-thick Au deposition in an e-beam evaporator. After fabrication of TMD devices, ~150 μm of pristine DNA, Ln-DNA, or Co-DNA sample solution was dropped onto TMD devices in order to fully cover the TMD surface. After drop-casting the sample on TMD, we incubated the samples in a natural drying environment at room temperature for about 10 hours. We performed this same drop-casting and natural drying process for each TMD transistor five times to ensure full coverage of DNA on the TMD flakes. After finishing this process, TMDs were washed with DI water. Transistors were doped by Ln-DNA or Co-DNA with different metal ions (Eu³⁺, Gd³⁺, Tb³⁺, Er³⁺, or Cu²⁺). Here, the thickness of the gate oxide (SiO₂) was 90 nm.

**Optical characterization of Ln-DNA or Co-DNA-doped TMD electronic devices.** In order to investigate the optoelectronic properties of the fabricated Ln-DNA- or Co-DNA-doped TMD devices, a current-voltage (I_D−I_C) measurement was performed under both dark and illuminated conditions. The light source was a diode laser with wavelengths of 655 nm, 785 nm, and 850 nm, and an optical power of 0.6 mW. For the characterization and comparison of the TMD optoelectronic devices doped by Ln-DNA or Co-DNA with different metal ions (Eu³⁺, Gd³⁺, Tb³⁺, Er³⁺, or Cu²⁺), photo-responsivity (R) and detectivity (D*) were calculated from I_D−I_C. Here, the effective area of the detector, e is the absolute value of electron charge (1.6 × 10⁻¹⁹ C), and I_Dark is the dark current.

**Photoluminescence measurement of pristine and Tb³⁺-doped DNA.** The emission spectra of pristine DNA and Tb³⁺-DNA on silica substrate were obtained using a fluorescence spectrophotometer (LS-55, PerkinElmer Instruments, USA) at room temperature. The emission spectrum was measured by exciting the sample at wavelength of 400 nm.

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Author Contributions
J.-H.P. and S.H.P. initiated the project. D.-H.K., S.R.D. and H.-Y.P. contributed to the experimental design and device fabrication. D.-H.K., H.-Y.P., J.S. and Y.R. analyzed the data. B.G. carried out the atomic force microscopy. J.J. and S.L. performed the Raman spectroscopy. S.R.D. carried out the photoluminescence measurement. D.-H.K. carried out electrical and optical characterization. All authors participated in writing the paper.

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