Polynomial Relations for $q$-Characters via the ODE/IM Correspondence

Juanjuan SUN

Graduate School of Mathematical Sciences, The University of Tokyo, Komaba, Tokyo 153-8914, Japan
E-mail: sunjuan@ms.u-tokyo.ac.jp

Received January 08, 2012, in final form May 10, 2012; Published online May 15, 2012
http://dx.doi.org/10.3842/SIGMA.2012.028

Abstract. Let $U_q(b)$ be the Borel subalgebra of a quantum affine algebra of type $X_n^{(1)}$ ($X = A, B, C, D$). Guided by the ODE/IM correspondence in quantum integrable models, we propose conjectural polynomial relations among the $q$-characters of certain representations of $U_q(b)$.

Key words: Borel subalgebra; $q$-character; Baxter’s $Q$-operator; ODE/IM correspondence

2010 Mathematics Subject Classification: 81R10; 17B37; 81R50

1 Introduction

Let $U_q(g)$ be a quantum affine algebra of type $X_n^{(1)}$ ($X = A, B, C, D$), and let $U_q(b)$ be its Borel subalgebra. In this paper we shall consider the problem of finding polynomial relations satisfied by the $q$-characters of the fundamental modules in the sense of [20] and related modules. This problem is intimately related with that of functional equations for Baxter’s $Q$-operators in quantum integrable models. In order to motivate the present study let us review this connection.

In quantum integrable systems, one is interested in the spectra of a commutative family of transfer matrices. The latter are constructed from the universal $R$ matrix of a quantum affine algebra, by taking the trace of the first component over some finite-dimensional representation called ‘auxiliary space’. When the auxiliary spaces are the Kirillov–Reshetikhin (KR) modules, the corresponding transfer matrices satisfy an important family of polynomial identities known as the $T$-system [24, 26]. (For a recent survey on this topic, see [25].) Subsequently the $T$-system has been formulated and proved [19, 28] as identities of $q$-characters. It has been shown further [21] that the $T$-system is actually the defining relations of the Grothendieck ring, which is a polynomial ring [17], of finite-dimensional modules of quantum affine algebras. Notice that, from the construction by trace, the $q$-characters and the transfer matrices are both ring homomorphisms defined on the Grothendieck ring. Since the $q$-characters are injective [16, 17], identities for $q$-characters imply the same identities for transfer matrices.

Baxter’s $Q$-operators were first introduced in the study of the 8-vertex model [1]. Since then they have been recognized as a key tool in classical and quantum integrable systems, and there is now a vast literature on this subject. In the seminal paper [4], Bazhanov, Lukyanov and Zamolodchikov revealed that the $Q$-operators can also be obtained from the universal $R$ matrix, provided the auxiliary space is chosen to be a (generically infinite-dimensional) representation of the Borel subalgebra. The work [4] for $U_q(\mathfrak{sl}_2)$ has been extended by several authors [2, 3, 5, 23, 29] to higher rank and supersymmetric cases.

In view of the results mentioned above, it is natural to ask whether one can find polynomial relations, analogous to the $T$-system, for the $q$-characters of $U_q(b)$ (see definitions in [20]) (and
hence for the $Q$-operators as well). The goal of this paper is to propose candidates of such identities.

Our idea is to use the so-called ODE/IM correspondence which relates the eigenvalues of $Q$-operators and certain ordinary linear differential equations. The reader is referred e.g.

to the review [12] on this topic. In [11], the correspondence is discussed for general non-twisted affine Lie algebras using scalar (pseudo-)differential operators. In this paper we reformulate their results in terms of first-order systems. The general setting (to be explained in Section 3) is as follows.

Let $\mathfrak{g}$ denote the Langlands dual Lie algebra of $\mathfrak{g}$. Consider the following $L_\mathfrak{g}$-valued first-order linear differential operator

$$\mathcal{L} = \frac{d}{dx} - \frac{\ell}{x} + \sum_{i=1}^{n} e_i + (x^{Mh^\vee} - E)e_0. \quad (1.1)$$

Here $e_0, \ldots, e_n$ are the Chevalley generators of $L_\mathfrak{g}$, $\ell$ is a generic element of the Cartan subalgebra $L_\mathfrak{h} \subset L_\mathfrak{g}$, $h^\vee$ is the dual Coxeter number of $\mathfrak{g}$, and $M > 0$, $E \in \mathbb{C}$ are parameters. For each fundamental representation $V^{(a)}$ of $L_\mathfrak{g}$, there is a basis $\{\chi^{(a)}_J(x, E)\}$ of $V^{(a)}$-valued solutions to the equation $\mathcal{L}\chi = 0$ characterized by the behavior

$$\chi^{(a)}_J(x, E) = u^{(a)}_J x^{\lambda^{(a)}_J} (1 + O(x)), \quad x \to 0,$$

where $\lambda^{(a)}_J$, $u^{(a)}_J$ are the eigenvalues of $\ell$ and the corresponding eigenvectors. Hence this basis is labeled by weight vectors of $V^{(a)}$. There is also a canonical solution $\psi^{(a)}(x, E)$ which has the fastest decay at $x \to +\infty$. Due to the special choice (1.1) of $\mathcal{L}$, the canonical solutions are shown to satisfy a system of relations similar to the Plücker relations, called the $\psi$-system in [11] (see Subsection 3.2 below). Introduce the connection coefficients $Q^{(a)}_J(E) \in \mathbb{C}$ by

$$\psi^{(a)}(x, E) = \sum_J Q^{(a)}_J(E) \chi^{(a)}_J(x, E).$$

Then the $\psi$-system implies a set of polynomial relations among the connection coefficients $Q^{(a)}_J(E)$.

We expect the following to be true:

(1) For each connection coefficient $Q^{(a)}_J(E)$, there is associated a formal power series $Q^{(a)}_{J,z}$. Up to a simple multiplicative factor, the latter is the $q$-character of an irreducible highest $\ell$-weight module of $U_q(\mathfrak{b})$. In particular, for the highest or lowest weights of $V^{(a)}$ the corresponding modules are the fundamental modules [20] of $U_q(\mathfrak{b})$.

(2) With the identification

$$Q^{(a)}_J(E) \leftrightarrow Q^{(a)}_{J,z}, \quad E \leftrightarrow z,$$

the polynomial relations for $Q^{(a)}_J(E)$ implied by the $\psi$-system hold true also for $Q^{(a)}_{J,z}$.

We adopt (1), (2) as our working hypothesis for finding relations among the $q$-characters. A few remarks are in order here. We consider separately the cases related to the spin representation (i.e., $(\mathfrak{g}, a) = (C_1^{(1)}, n), (D_1^{(1)}, n-1), (D_1^{(1)}, n)$) and use the letter $R_{\varepsilon,z}^{(a)}$ to denote the counterpart of $Q^{(a)}_{J,z}$. We expect that $Q^{(a)}_{J,z}$ or $R_{\varepsilon,z}^{(a)}$ corresponds to the $q$-character of an irreducible highest $\ell$-weight module of $U_q(\mathfrak{b})$ whose highest $\ell$-weight is given by formulas (4.16), (4.19)–(4.20), (4.23)–(4.24). (In the case $C_1^{(1)}$, there are problems with this interpretation, as
explained in Remark 5.9.) Also the relations for the \(Q_{J,z}^{(a)}\)'s or \(R_{\epsilon,z}^{(a)}\)'s are not exactly the same as those for \(Q_{J}^{(a)}(E)\)'s, but it is necessary to fine-tune the coefficients by some power functions independent of \(z\). The details will be given in Section 5 below. That it is natural to consider all \(Q_{J,z}^{(a)}\) corresponding to general weights of \(V^{(a)}\) is a viewpoint suggested by the work [29] for type \(A\) algebras.

Now let us come to the content of the present work. As a first step, we give explicit candidates for \(Q_{i,z}^{(1)}\) associated with each weight of the vector representation \(V^{(1)}\). This is done by taking suitable limits of the known \(q\)-characters of KR modules given by tableaux sums. It is known [20] that for the highest and lowest weights this procedure indeed gives the irreducible \(q\)-characters. As the next step, we define the formal series \(Q_{J,z}^{(a)}\) for a general node \(a\) of the Dynkin diagram. We define them by Casorati determinants whose entries are \(Q_{i,z}^{(1)}\) with suitable shifts of parameters.

We then give candidates of polynomial relations among the \(Q_{i,z}^{(a)}\) expected from the \(\psi\)-system. In the cases related to spin representations, however, we do not have explicit candidates the \(R_{\epsilon,z}^{(a)}\)'s in general, so we only write down the candidates for the relations. The resulting polynomial relations are given in Proposition 5.1, Theorem 5.2 (type \(A_{n}^{(1)}\)), Conjecture 5.3 (type \(B_{n}^{(1)}\)), Conjectures 5.5–5.8 (type \(C_{n}^{(1)}\)) and Conjectures 5.10, 5.11 (type \(D_{n}^{(1)}\), respectively.

For the type \(A\) algebras, one can check that \(Q_{i,z}^{(1)}\) (not necessarily the highest or lowest ones) are indeed irreducible \(q\)-characters of modules given in [23]. The polynomial relations corresponding to the \(\psi\)-system can be summarized as a single identity

\[
\det \left( Q_{\nu,q}^{(1)} \right)_{\mu,\nu=1}^{n+1} = 1, \tag{1.2}
\]

where \(x_{\nu} = e^{\nu}, \{e_{\nu}\}_{\nu=1}^{n+1}\) being an orthonormal basis related to the simple roots by \(e_{\nu} = \epsilon_{\nu} - \epsilon_{\nu+1}\). This relation (as an identity for \(Q\)-operators) has been known to experts under the name ‘Wronskian identity’. We shall give a direct proof that (1.2) is satisfied for \(q\)-characters in Appendix B.

For the other types of algebras, the situation is less satisfactory. At the moment we do not know the irreducibility of modules corresponding to the \(Q_{i,z}^{(1)}\) given by our procedure (except those corresponding to the highest or lowest vectors). For the spin representations of \(C_{n}^{(1)}\) and \(D_{n}^{(1)}\), explicit formulas for the \(R_{\epsilon,z}^{(a)}\)'s are missing. More seriously, we have not been able to prove the proposed identities for \(q\)-characters by computational methods. Instead, we support our working hypothesis by performing the following checks:

1. \(\mathfrak{g} = B_{2}^{(1)}\), proof by hand,
2. \(\mathfrak{g} = B_{3}^{(1)}\), computer check,
3. \(\mathfrak{g} = D_{4}^{(1)}\), computer check,
4. \(\mathfrak{g} = B_{4}^{(1)}, C_{4}^{(1)}, D_{4}^{(1)}\), proof in the limit to ordinary characters.

The main results of the present paper consist in formulating the conjectured relations, and in performing the checks mentioned above.

The text is organized as follows. In Section 2, we prepare basic definitions concerning the Borel subalgebra \(U_q(\mathfrak{b})\) of a quantum affine algebra \(U_q(\mathfrak{g})\), and collect necessary facts about their representations. In Section 3, we give an account of the \(\psi\)-system in the ODE/IM correspondence and indicate how to derive them by using the formulation by first-order systems. We note that in [11] the \(\psi\)-system for algebras other than \(A\) type is mentioned as conjectures. In Section 4, we introduce the series \(Q_{i,z}^{(1)}\) as limits of the \(q\)-characters of KR modules. We also define these series for the other nodes of the Dynkin diagram. Section 5 is devoted to the proposals for polynomial
relations. By comparing with the relations for the connection coefficients, we write down the relations for each type of algebras $A_n^{(1)}$, $B_n^{(1)}$, $C_n^{(1)}$ and $D_n^{(1)}$. In Section 6, we give a summary of our work.

The text is followed by four appendices. Appendix A gives a list of realizations of the dual Lie algebras $\hat{\mathfrak{g}}$. In Appendix B we give a proof of the Wronskian identity for type $A_n^{(1)}$. In Appendix C, we prove the identities for type $B_n^{(1)}$ in the limit to ordinary characters. In Appendix D the same is done for type $C_n^{(1)}$ and $D_n^{(1)}$.

2 Preliminaries

In this section we introduce our notation on quantum affine algebras and their Borel subalgebras, and collect necessary facts that will be used later. Throughout this paper, we assume that $q$ is a nonzero complex number which is not a root of unity.

2.1 Quantum Borel algebras

Let $\mathfrak{g}$ be an affine Lie algebra associated with a generalized Cartan matrix $C = (c_{ij})_{0 \leq i,j \leq n}$ of non-twisted type. Let $D = \text{diag}(d_0, \ldots, d_n)$ be the unique diagonal matrix such that $DC$ is symmetric and $d_0 = 1$. Set $I = \{1,2,\ldots,n\}$, and let $\check{\mathfrak{g}}$ denote the simple Lie algebra with the Cartan matrix $(c_{ij})\ i, j \in I$. Let $\{\alpha_i\}_{i \in I}$, $\{\alpha_i^\vee\}_{i \in I}$ and $\{\omega_i\}_{i \in I}$ be the simple roots, simple coroots and the fundamental weights of $\check{\mathfrak{g}}$, respectively. We set $P = \bigoplus_{i \in I}\mathbb{Z}\omega_i$, $Q = \bigoplus_{i \in I}\mathbb{Z}\alpha_i$.

Set $q_i = q^{d_i}$. We shall use the standard notation

$$[k]_i = \frac{q^k_i - q^{-k}_i}{q_i - q_i^{-1}}, \quad [n]_i = \prod_{k=1}^{n}[k]_i, \quad \left[\begin{array}{c} n \\ k \end{array}\right]_i = \frac{[n]_i!}{[k]_i!\ [n-k]_i!}.$$ 

The quantum affine algebra $U_q(\mathfrak{g})$ is the $\mathbb{C}$-algebra defined by generators $E_i$, $F_i$, $K_i^{\pm 1}$ $(i = 0, \ldots, n)$ and the relations

$$K_i K_i^{-1} = 1 = K_i^{-1} K_i, \quad K_i K_j = K_j K_i, \quad K_i E_j K_i^{-1} = q^{c_{ij}}_i E_j, \quad K_i F_j K_i^{-1} = q^{-c_{ij}}_i F_j, \quad [E_i, F_j] = \delta_{ij} \frac{K_i - K_i^{-1}}{q_i - q_i^{-1}},$$

$$\sum_{r=0}^{1-c_{ij}} \left[\begin{array}{c} 1 - c_{ij} \\ r \end{array}\right]_i (-1)^r E_i^r E_j E_i^{1-c_{ij}-r} = 0, \quad i \neq j,$$

$$\sum_{r=0}^{1-c_{ij}} \left[\begin{array}{c} 1 - c_{ij} \\ r \end{array}\right]_i (-1)^r F_i^r F_j F_i^{1-c_{ij}-r} = 0, \quad i \neq j.$$ 

We do not write the formulas defining the Hopf structure on $U_q(\mathfrak{g})$ since we are not going to use them.

As is well known [6, 13], $U_q(\mathfrak{g})$ is isomorphic to the $\mathbb{C}$-algebra with generators $x_{i,r}^\pm$ $(i \in I, \ r \in \mathbb{Z})$, $k_i^{\pm 1}$ $(i \in I)$, $h_{i,r}$ $(i \in I, \ r \in \mathbb{Z}\setminus\{0\})$ and central elements $c^{\pm 1/2}$, with the following defining relations

$$k_i k_i^{-1} = 1 = k_i^{-1} k_i, \quad c^{1/2} c^{-1/2} = 1 = c^{-1/2} c^{1/2}, \quad k_i k_j = k_j k_i, \quad k_i h_{j,r} = h_{j,r} k_i,$$

$$k_i x_{j,r}^\pm k_i^{-1} = q_i^{c_{ij}} x_{j,r}^\pm, \quad [h_{i,r}, x_{j,s}^\pm] = \pm \frac{1}{r} [r c_{ij}] c^{r|\pm|/2} x_{j,r+s}^\pm,$$

$$x_{i,r+1}^\pm x_{j,s}^\pm - q_i^{c_{ij}} x_{j,s}^\pm x_{i,r+1}^\pm = q_i^{c_{ij}} x_{i,r}^\pm x_{j,s+1}^\pm - x_{j,s+1}^\pm x_{i,r}^\pm.$$
Polynomial Relations for $q$-Characters via the ODE/IM Correspondence

\[ [x_{i,r}^+, x_{j,s}^-] = \delta_{ij} \frac{c^{(r-s)/2} \phi_{i,r+s}^+ - c^{-(r-s)/2} \phi_{i,r+s}^-}{q_i - q_i^{-1}}, \]

\[ \sum_{\pi \in \mathfrak{S}_{1-c_{ij}}} \sum_{k=0}^{1-c_{ij}} (-1)^k \left[ 1 - c_{ij} \right] x_{i,r(1)}^+ \cdots x_{i,r(k)}^+ x_{j,s}^+ x_{i,r(k+1)}^+ \cdots x_{i,r(1-c_{ij})}^+ = 0, \quad i \neq j, \]

for all integers $r_j$, where $\mathfrak{S}_m$ is the symmetric group on $m$ letters, and the $\phi_{i,r}^\pm$ are given by

\[ \sum_{r=0}^{\infty} \phi_{i,r}^\pm u^{\pm r} = k_i^{\pm 1} \exp \left( \pm (q_i - q_i^{-1}) \sum_{s=1}^{\infty} h_{i,s} u^{\pm s} \right). \]

By definition, the Borel subalgebra $U_q(b)$ is the Hopf subalgebra of $U_q(g)$ generated by $E_i$, $K_i^{\pm 1}$ $(0 \leq i \leq n)$. It is known [22] that $U_q(b)$ is isomorphic to the algebra with generators $E_i$, $K_i^{\pm 1}$ $(0 \leq i \leq n)$ and the defining relations

\[ K_i K_j = K_j K_i, \quad K_i E_j K_i^{-1} = q_i^{c_{ij}} E_j, \]

\[ \sum_{r=0}^{1-c_{ij}} \left[ 1 - c_{ij} \right] (-1)^r E_i^r E_j E_i^{1-c_{ij}-r} = 0, \quad i \neq j. \]

The algebra $U_q(b)$ contains $x_{i,m}^+$, $x_{i,r}^-$, $k_i^{\pm 1}$ and $\phi_{i,r}^\pm$, where $i \in I$, $m \geq 0$ and $r > 0$.

### 2.2 Category $\mathcal{O}$ of $U_q(b)$

In this subsection we recall basics about $U_q(b)$-modules in category $\mathcal{O}$. For more details see [20].

Denote by $\mathfrak{t}_\ell$ the subalgebra of $U_q(b)$ generated by $\{k_i^{\pm 1}\}_{i \in I}$. For a $U_q(b)$-module $V$ and $\lambda \in \mathfrak{p}$, set

\[ V_\lambda = \{ v \in V \mid k_i v = q_i^{\lambda(\alpha_i^\vee)} v, \ \forall i \in I \}. \]

When $V_\lambda \neq 0$, it is called the weight space of weight $\lambda$. A $U_q(b)$-module $V$ is said to be of type 1 if $V = \bigoplus_{\lambda \in \mathfrak{p}} V_\lambda$.

A series of complex numbers $\Psi = (\Psi_{i,r})_{i \in I, r \in \mathbb{Z}_{\geq 0}}$ is called an $\ell$-weight if $\Psi_{i,0} \neq 0$ for all $i \in I$. We denote by $\mathfrak{t}_\ell^*$ the set of $\ell$-weights. For a $U_q(b)$-module $V$ and $\Psi \in \mathfrak{t}_\ell^*$, the subspace

\[ V(\Psi) = \{ v \in V \mid \exists p > 0, \ \forall i \in I, \ \forall m \geq 0, (\phi_{i,m}^+ - \Psi_{i,m})^p v = 0 \} \]

is called the $\ell$-weight space of $\ell$-weight $\Psi$.

A $U_q(b)$-module $V$ is said to be a highest $\ell$-weight module of highest $\ell$-weight $\Psi \in \mathfrak{t}_\ell^*$ if there exists a vector $v \in V$ such that $V = U_q(b) v$ and

\[ E_i v = 0, \quad i \in I, \quad \phi_{i,r}^+ v = \Psi_{i,r} v, \quad i \in I, \ r \geq 0. \]

For each $\Psi \in \mathfrak{t}_\ell^*$ there exists a unique simple $U_q(b)$-module of highest $\ell$-weight $\Psi$. We denote it by $L(\Psi)$.

A highest $\ell$-weight module is of type 1 if its highest $\ell$-weight $\Psi$ satisfies

\[ \Psi_{i,0} = q_i^{p_i} \quad \text{for some } p_i \in \mathbb{Z}, \ i \in I. \quad (2.1) \]

For any non-zero complex numbers $c_i \in \mathbb{C}^\times$, the map $E_i \mapsto E_i$, $K_i \mapsto c_i K_i$ $(i = 0, \ldots, n)$ gives rise to an automorphism of $U_q(b)$. After twisting by such an automorphism, any highest $\ell$-weight module can be brought to one satisfying the condition (2.1). We denote by $\mathfrak{t}_{\ell,p}^*$ the set of $\ell$-weights satisfying (2.1).

Set $D(\lambda) = \lambda - Q^+, \ Q^+ = \sum_{i \in I} \mathbb{Z}_{\geq 0} \alpha_i$. A $U_q(b)$-module $V$ of type 1 is said to be an object in category $\mathcal{O}$ if
1. for all $\lambda \in P$ we have $\dim V_\lambda < \infty$,
2. there exist a finite number of elements $\lambda_1, \ldots, \lambda_s \in P$ such that the weights of $V$ are contained in $\bigcup_{j=1, \ldots, s} D(\lambda_j)$.

In what follows we shall identify $\Psi \in t_i^*$ with their generating series,

$$\Psi = (\Psi_1(u), \ldots, \Psi_n(u)), \quad \Psi_i(u) = \sum_{r \geq 0} \Psi_{i,r} u^r.$$ 

Simple objects in category $O$ are classified by the following theorem.

**Theorem 2.1** ([20]). Suppose that $\Psi \in t_i^*$. Then the simple module $L(\Psi)$ is an object in category $O$ if and only if $\Psi_i(u)$ is a rational function of $u$ for any $i \in I$.

In particular, for $i \in I$ and $z \in \mathbb{C}^x$, the simple modules $L_{i,z} = L(\Psi)$ defined by the highest $\ell$-weight

$$\Psi_j(u) = \begin{cases} (1 - zu)\pm 1, & j = i, \\ 1, & j \neq i, \end{cases}$$

are objects in category $O$. These modules are called the fundamental representations [20].

It is known [7, 8] that finite-dimensional simple $U_q(\mathfrak{g})$-modules remain simple when restricted to $U_q(\mathfrak{b})$. According to the classification of the former [9, 10], the simple module $L(\Psi)$ is finite-dimensional if its highest $\ell$-weight has the form

$$\Psi_i(u) = q_i^{\deg P_i} \frac{P_i(q_i^{-1}u)}{P_i(q_i u)}, \quad \forall i \in I,$$

where $P_i(u)$ is a polynomial such that $P_i(0) = 1$. In the case where

$$P_j(u) = \begin{cases} (1 - q_i^{-m+1}zu)(1 - q_i^{-m+3}zu) \cdots (1 - q_i^{-m-1}zu), & j = i, \\ 1, & j \neq i, \end{cases}$$

with some $i \in I$, $m \in \mathbb{Z}_{>0}$ and $z \in \mathbb{C}^x$, the module $L(\Psi)$ is called a Kirillov–Reshetikhin (KR) module. We denote it by $W_{m,z}^{(i)}$.

### 2.3 Characters and $q$-characters

We recall the definition of $q$-characters (see [20]) and characters of representations of $U_q(\mathfrak{b})$. Let $Z_{i}^{U,P}$ denote the set of maps $t_i^* \rightarrow \mathbb{Z}$. For $\Psi \in t_i^*$, define $[\Psi] \in Z_{i}^{U,P}$ by $[\Psi](\Psi') = \delta_{\Psi,\Psi'}$. For a $U_q(\mathfrak{b})$-module $V$ of type 1 in category $O$, its $q$-character $\chi_q(V)$ is defined as an element of $Z_{i}^{U,P}$,

$$\chi_q(V) = \sum_{\Psi \in t_i^*} \dim V(\Psi) \cdot [\Psi].$$

Similarly let $Z^P$ denote the set of maps $P \rightarrow \mathbb{Z}$, and define $e^\lambda \in Z^P$ by $e^\lambda(\mu) = \delta_{\lambda,\mu}$. The ordinary character $\chi(V)$ is an element of $Z^P$,

$$\chi(V) = \sum_{\lambda \in P} \dim V_\lambda \cdot e^\lambda.$$

We have a natural map $\varpi : Z_{i}^{U,P} \rightarrow Z^P$ which sends $[\Psi]$ to $e^\lambda$ such that $\Psi_{i,0} = q_i^{\lambda(\alpha_i^\vee)}$. Under $\varpi$ the $q$-character specializes to the ordinary character,

$$\chi(V) = \varpi(\chi_q(V)).$$
3 $\psi$-system

In this section, we reformulate the $\psi$-systems given in [11].

3.1 $Lg$-connection

From now on, let $g$ be an affine Lie algebra of type $X_n^{(1)} (X = A, B, C, D)$, and let $Lg$ denote its Langlands dual algebra. Let $h^\vee$ be the dual Coxeter number of $g$ (see Table 1).

|   | $A_n^{(1)}$ | $B_n^{(1)}$ | $C_n^{(1)}$ | $D_n^{(1)}$ |
|---|-------------|-------------|-------------|-------------|
| $Lg$ | $A_n^{(1)}$ | $A_{2n-1}^{(2)}$ | $D_{n+1}^{(2)}$ | $D_n^{(1)}$ |
| $h^\vee$ | $n + 1$ | $2n - 1$ | $n + 1$ | $2n - 2$ |
| $\dim V^{(1)}$ | $n + 1$ | $2n$ | $2n + 2$ | $2n$ |

Denote by $e_j, f_j, h_j$ ($0 \leq j \leq n$) the Chevalley generators of $Lg$. We set $e = \sum_{j=1}^n e_j$. Fix also an element $\ell \in Lh$ from the Cartan subalgebra of $Lg$, and let $\xi \in \mathbb{C}^\times$. We consider the following $Lg$-valued connection (cf. e.g. [15]):

$$L = \frac{d}{dx} - \frac{\ell}{x} + e + p(x, E)\xi e_0, \quad (3.1)$$

$$p(x, E) = x^{Mh^\vee} - E, \quad M > 0, \quad E \in \mathbb{C}. \quad (3.2)$$

Take an element $h_\rho \in Lh$ such that $[h_\rho, e_j] = e_j$ ($1 \leq j \leq n$) and $[h_\rho, e_0] = -(h^\vee - 1)e_0$. The choice (3.2) ensures the following symmetry property for $L = L(x, E; \xi)$:

$$\omega^{k(h_\rho - 1)} L(x, E; e^{2\pi i k} \xi) \omega^{-kh_\rho} = L(\omega^k x, \Omega^k E; \xi), \quad (3.3)$$

where $k \in \mathbb{C}$ and

$$\omega = e^{\frac{2\pi i k}{M+1}}, \quad \Omega = \omega^{h^\vee M}. \quad (3.4)$$

On any finite-dimensional $Lg$-module $V$, (3.1) defines a first-order system of differential equations $L(x, E; 1)\phi(x, E) = 0$. Quite generally, for a $Lg$-module $V$, we denote by $V_k$ the $Lg$-module obtained by twisting $V$ by the automorphism $e_j \mapsto \exp(2\pi i k \delta_{j,0}) e_j, f_j \mapsto \exp(-2\pi i k \delta_{j,0}) f_j$. The operator $L(x, E; e^{2\pi i k})$ represents the action of $L(x, E; 1)$ on $V_k$. For a $V$-valued solution $\phi(x, E)$ we set

$$\phi_k(x, E) = \omega^{-kh_\rho} \phi(\omega^k x, \Omega^k E). \quad (3.5)$$

Then the symmetry (3.3) implies $L(x, E; e^{2\pi i k}) \phi_k(x, E) = 0$.

With each node $a$ of the Dynkin diagram of $Lg$ is associated a fundamental module $V^{(a)}$ of $Lg$. We summarize our convention about them and some facts which will be used later. We leave the proofs to the readers (see Subsection 5.3 for an example).

Remark 3.1. In [11], scalar (pseudo-)differential equations are considered. Using the realization of $Lg$ given in Appendix A and rewriting the equation $L\phi = 0$ for the highest component of $\phi$, one obtains the formulas [11, (3.18)–(3.21)] (for simplicity we have taken $K = 1$ there).
The module \( V^{(1)} \) is called the vector representation of \( L \mathfrak{g} \). Its explicit realization is given in Appendix A. If \( k \) is an integer, it is obvious that \( V_k = V \) for any \( V \). The vector representation for \( \mathfrak{g} = C_n^{(1)} \) has the additional property
\[
V^{(1)}_\frac{1}{2} \cong V^{(1)}, \quad \mathfrak{g} = C_n^{(1)}.
\] (3.6)

For general \( a \), we distinguish the following two cases,
\[
\text{(NS)}: \quad (\mathfrak{g}, a) \neq (C_n^{(1)}, n), (D_n^{(1)}, n-1), (D_n^{(1)}, n),
\]
\[
\text{(S)}: \quad (\mathfrak{g}, a) = (C_n^{(1)}, n), (D_n^{(1)}, n-1), (D_n^{(1)}, n).
\]
We shall refer to them as the non-spin case and the spin case, respectively. Here and after we set
\[
t = \begin{cases} 
2, & \mathfrak{g} = C_n^{(1)}, \\
1, & \text{otherwise}.
\end{cases}
\]
In the non-spin case, we have
\[
V^{(a)} = \bigwedge_a V^{(1)}_{\frac{1}{2} + \frac{a}{2n}} \quad \text{for (NS).}
\] (3.7)

In the case \( \mathfrak{g} = B_n^{(1)} \), we have in addition
\[
V^{(a)}_\frac{1}{2} \cong V^{(2n-a)}, \quad \mathfrak{g} = B_n^{(1)}, \quad a = 1, \ldots, n,
\] (3.8)

where \( V^{(a)} \) for \( a > n \) stands for the right-hand side of (3.7).

In the case \( (\mathfrak{g}, a) = (C_n^{(1)}, n) \), \( V^{(n)} \) is the spin representation of the subalgebra \( L \hat{\mathfrak{g}} = \mathfrak{o}(2n+1) \subseteq D_n^{(2)} \). Likewise, in the cases \( (\mathfrak{g}, a) = (D_n^{(1)}, n-1) \) or \( (D_n^{(1)}, n) \), \( V^{(n-1)} \) and \( V^{(n)} \) are the spin representations of the subalgebra \( L \hat{\mathfrak{g}} = \mathfrak{o}(2n) \subseteq D_n^{(1)} \).

Let us consider the solutions at the irregular singularity \( x = \infty \). It is convenient to use a gauge transformed form of \( \mathcal{L} \),
\[
x^{Mh_\rho} \mathcal{L} x^{-Mh_\rho} = \frac{d}{dx} + \Lambda x^M - \ell + Mh_\rho - E_0 x^{-M(h_\rho - 1)},
\] (3.9)

where \( \Lambda = e + E_0 \). Let \( \mu^{(a)} \) be the eigenvalue of \( \Lambda \) on \( V^{(a)} \) which has the largest real part. This eigenvalue is multiplicity free, and is given explicitly as follows:
\[
\mu^{(a)} = \begin{cases} 
\sin \frac{\pi a}{2h_\rho}, & \text{for (NS)}, \\
\sin \frac{\pi a}{2h_\rho} + \frac{1}{2} \sin \frac{\pi a}{2h_\rho}, & \text{for (S)}.
\end{cases}
\]

Let \( \mathbf{u}^{(a)} \) be an eigenvector of \( \Lambda \) corresponding to \( \mu^{(a)} \). From the representation (3.9) it follows that there is a unique \( V^{(a)} \)-valued solution \( \psi^{(a)}(x, E) \) which satisfies the following in a sector containing the positive real axis \( x > 0 \):
\[
\psi^{(a)}(x, E) = e^{-\frac{\mu^{(a)}}{Mh_\rho} x^{M+1}} x^{-Mh_\rho} \left( \mathbf{u}^{(a)} + o(1) \right) \quad (x \to \infty).
\]

We call \( \psi^{(a)}(x, E) \) the canonical solution. In view of the relation (3.7) and the formula for \( \mu^{(a)} \) given above, we have
\[
\psi^{(a)} = \psi^{(1)}_{\frac{a+1}{2n}} \wedge \psi^{(1)}_{\frac{a+1}{2n}} \wedge \cdots \wedge \psi^{(1)}_{\frac{a-1}{2n}} \quad \text{for (NS).}
\] (3.10)

Here \( \psi^{(i)} \) is defined by (3.5).
3.2 \( \psi \)-system

In this subsection we state the \( \psi \)-system for the canonical solutions \( \psi^{(a)}(x, E) \) introduced above. Let us consider them case by case.

\[
g = A_n^{(1)} \quad (Lg = A_n^{(1)}).
\]

For \( a = 1, \ldots, n \), we have the embedding of \( Lg \)-modules

\[
\iota : \bigwedge^2 V_1^{(a)} \hookrightarrow V^{(a-1)} \otimes V^{(a+1)},
\]

(3.11)

where \( V^{(0)} = V^{(n+1)} = \mathbb{C} \). The explicit expression of the embedding is given in Subsection 5.1. On the space \( V^{(a-1)} \otimes V^{(a+1)} \), the functions \( \phi = \iota(\psi^{(a)}_{1/2} \wedge \psi^{(a)}_{1/2}) \) and \( \phi = \psi^{(a-1)} \otimes \psi^{(a+1)} \) both satisfy the equation \( \mathcal{L}\phi = 0 \) and have the behavior

\[
\phi(x, E) = O \left( \exp \left( -\frac{2\mu(a) \cos \frac{\pi}{M+1} x^{M+1}}{x} \right) \right), \quad x \to \infty
\]

in a sector containing \( x > 0 \). Since such a solution is unique up to a constant multiple, we conclude that (after adjusting the constant multiple)

\[
\iota(\psi^{(a)}_{1/2} \wedge \psi^{(a)}_{1/2}) = \psi^{(a-1)} \otimes \psi^{(a+1)}, \quad a = 1, \ldots, n.
\]

(3.12)

In particular, denoting by \( \{u_j\}^{n+1}_{j=1} \) the standard basis of \( V^{(1)} \) we have

\[
\psi^{(n+1)} = u_1 \wedge \cdots \wedge u_{n+1}.
\]

(3.13)

We call the relations (3.12), (3.13) the \( \psi \)-system for \( A_n^{(1)} \).

The \( \psi \)-system for the other types can be deduced by the same argument, using the relevant embeddings of representations. We obtain relations of the following form.

\[
g = B_0^{(1)} \quad (Lg = A_{2n-1}^{(2)}).
\]

\[
\iota(\psi^{(a)}_{1/2} \wedge \psi^{(a)}_{1/2}) = \psi^{(a-1)} \otimes \psi^{(a+1)}, \quad a = 1, \ldots, n-1,
\]

(3.14)

\[
\iota(\psi^{(a)}_{1/4} \wedge \psi^{(a)}_{1/4}) = \psi^{(a-1)}_{-1/4} \otimes \psi^{(a+1)}_{1/4},
\]

(3.15)

Here \( \iota \) stands for the embedding of \( Lg \) (3.11) or

\[
\bigwedge^2 V_{1/4}^{(n)} \hookrightarrow V_{-1/4}^{(n-1)} \otimes V_{1/4}^{(n-1)},
\]

which follows from (3.11) and (3.8).

\[
g = C_0^{(1)} \quad (Lg = D_{n+1}^{(2)}).
\]

\[
\iota(\psi^{(a)}_{1/4} \wedge \psi^{(a)}_{1/4}) = \psi^{(a-1)}_{-1/4} \otimes \psi^{(a+1)}_{1/4}, \quad a = 1, \ldots, n-2,
\]

(3.16)

\[
\iota(\psi^{(a)}_{-1/2} \wedge \psi^{(a)}_{1/2}) = \psi^{(a-1)}_{-1/2} \wedge \psi^{(a+1)}_{1/2},
\]

(3.17)

\[
\iota(\psi^{(a)}_{-1/4} \wedge \psi^{(a)}_{1/4}) = \psi^{(a-1)}_{-1/4} \otimes \psi^{(a+1)}_{1/4},
\]

(3.18)

Here we have set

\[
\phi^{(n-1)} = \psi^{(1)}_{-n-2} \wedge \psi^{(1)}_{-n-4} \wedge \cdots \wedge \psi^{(1)}_{-n-2}, \quad \phi^{(n)} = \psi^{(1)}_{-n-1} \wedge \psi^{(1)}_{-n-3} \wedge \cdots \wedge \psi^{(1)}_{-n-2},
\]
and \( \iota \) stands for an analog of (3.11) or the embeddings
\[
\bigwedge_{n-2} V^{(1)}_{n-2} \simeq V^{(n-1)} \hookrightarrow V^{(n)}_{\frac{n}{2}}, \quad \bigwedge_{n-1} V^{(1)}_{n-1} \hookrightarrow V^{(n)}_{\frac{n-1}{2}} \otimes V^{(n)}_{\frac{n}{2}},
\]
where (3.6) is taken into account.

\[
\varrho = D^{(1)}_{n-1} (L \varrho = D^{(1)}_{n}).
\]

\[
\iota(\psi^{(a)}_{\frac{n}{2}} \wedge \psi^{(a)}_{\frac{n}{2}}) = \psi^{(a-1)} \otimes \psi^{(a+1)}, \quad a = 1, \ldots, n-3,
\]

(3.19)

\[
\iota(\psi^{(n-2)}_{\frac{n}{2}}) = \psi^{(n-1)}_{\frac{n}{2}} \wedge \psi^{(n-1)}_{\frac{n}{2}}, \quad \iota(\psi^{(n-2)}) = \psi^{(n)}_{\frac{n}{2}} \wedge \psi^{(n)}_{\frac{n}{2}},
\]

(3.20)

\[
\iota(\phi^{(n-1)}) = \psi^{(n-1)} \otimes \psi^{(n)}.
\]

(3.21)

We have set
\[
\phi^{(n-1)} = \psi^{(1)}_{\frac{n}{2}} \wedge \psi^{(1)}_{\frac{n}{2}} \wedge \cdots \wedge \psi^{(1)}_{\frac{n}{2}}
\]
and \( \iota \) stands for the embedding (3.11) or
\[
\bigwedge_{n-2} V^{(1)}_{n-2} \simeq V^{(n-2)} \hookrightarrow \bigwedge_{n-1} V^{(n-1)} \hookrightarrow \bigwedge_{n} V^{(n)}.
\]

3.3 Connection coefficients

Now we introduce the connection coefficients \( Q^{(n)}(E) \). First let us consider the vector representation \( V^{(1)} \). We choose \( \ell \) generic. Set \( \ell = \text{diag}(\ell_1, \ldots, \ell_N) \) \((N = \dim V^{(1)}, \text{see Appendix A})\), and let \( u_j \) \((1 \leq j \leq N)\) be the corresponding eigenvector of \( \ell \). Then there is a unique \( V^{(1)} \)-valued solution \( \chi^{(1)}_{j}(x, E) \) characterized by the expansion at the origin,
\[
\chi^{(1)}_{j}(x, E) = x^\ell u_j (1 + O(x)), \quad x \to 0.
\]

From the symmetry (3.3) of \( \mathcal{L}(x, E) \) we find that
\[
\chi^{(1)}_{j,k}(x, E) = \omega^{k \lambda_j} \chi^{(1)}_{j}(x, E), \quad k \in (1/t)\mathbb{Z},
\]

(3.22)

where we have set
\[
\ell - h_\rho = \text{diag}(\lambda_1, \ldots, \lambda_N).
\]

(3.23)

Define \( Q^{(1)}_{j}(E) \in \mathbb{C} \) by
\[
\psi^{(1)}(x, E) = \sum_{j=1}^{N} Q^{(1)}_{j}(E) \chi^{(1)}_{j}(x, E).
\]

(3.24)

From (3.22) and (3.24), we have for \( k \in (1/t)\mathbb{Z} \)
\[
\psi^{(1)}_{k}(x, E) = \sum_{j=1}^{N} Q^{(1)}_{j,k}(E) \omega^{k \lambda_j} \chi^{(1)}_{j}(x, E),
\]
Then the limiting procedure

Below we shall use the following elements of irreducible $U_q$-characters of the Borel subalgebra arise as certain limits of the KR modules. In this section we follow the same procedure to obtain a family of formal power series $Q^{(1)}_{i,z}$ associated with each weight space of the vector representation of $L^+_b$. Up to simple overall multipliers, those corresponding to the highest or lowest weights are the irreducible $q$-characters $\chi_q(L^+_b)$. We expect that in general the $Q^{(1)}_{i,z}$ are also proportional to irreducible $q$-characters of $U_q(b)$. (See however Remark 5.9.)

4 Series $Q^{(a)}_{i,z}$, $R^{(a)}_{\varepsilon,z}$

It has been shown in [20] that the fundamental modules $L^\pm_{i,z}$ of the Borel subalgebra arise as certain limits of the KR modules. In this section we follow the same procedure to obtain a family of formal power series $Q^{(1)}_{i,z}$ associated with each weight space of the vector representation of $L^+_b$. Up to simple overall multipliers, those corresponding to the highest or lowest weights are the irreducible $q$-characters $\chi_q(L^+_b)$. We expect that in general the $Q^{(1)}_{i,z}$ are also proportional to irreducible $q$-characters of $U_q(b)$. (See however Remark 5.9.)

Recall in Subsection 2.3, for $\Psi \in L^\ell_{\varepsilon,p}$, the element $[\Psi] \in \mathbb{Z}^\ell_{\varepsilon,p}$ is defined by $[\Psi]([\Psi]) = \delta_{\Psi,\Psi'}$. Below we shall use the following elements of $\mathbb{Z}^\ell_{\varepsilon,p}$:

$$Y_{i,z} = [(1, \ldots, (1 - zu)^{-1}, \ldots, 1), \quad e^{\omega_i} = [(1, \ldots, q_i, \ldots, 1)],$$

$$Y_{i,z} = [(1, \ldots, 1 - q_i^{-1}zu, \ldots, 1)] = e^{\omega_i} Y_{i,q_i} Y_{i,q_i^{-1}z},$$

$$A_{i,z} = Y_{i,q_i^{-1}z} Y_{i,q_i} \prod_{1 \leq j \leq n, c_{ij} = -1} Y_{j,z}^{-1} \prod_{1 \leq j \leq n, c_{ij} = -2} Y_{j,q_j^{-1}z} Y_{j,q_j}^{-1} \prod_{1 \leq j \leq n, c_{ij} = -3} Y_{j,q_j^{-1}z} Y_{j,q_j}^{-1} Y_{j,q_j}^{-1}.$$  

Highest $\ell$-weights are monomials in $Y_{i,z}^{\pm 1}$ and $e^{\pm \omega_i}$. Abusing the notation, for a monomial $M = [\Psi]$ we shall also write $L(M)$ for $L(\Psi)$.

4.1 The limiting procedure

Let us illustrate on examples the procedure for taking the limit.

Example 4.1 ($g = A_2^{(1)}$). We consider first the case $g = A_2^{(1)}$. Following [26, 27], we write

$$[1]_z = Y_{1,z}, \quad [2]_z = Y_{1,q_2}^{-1} Y_{2,q_2}, \quad [3]_z = Y_{2,q_3}^{-1}.$$ 

Then the $q$-character $\chi_q(W^{(1)}_{m,z})$ of the KR module $W^{(1)}_{m,z}$ is presented as a sum over the tableaux

$$\sum_{k_1, k_2, k_3 \geq 0 \atop k_1 + k_2 + k_3 = m} \frac{k_1 \cdot 1 \cdot 2 \cdot 3 \cdots k_1 k_2 k_3}{k_1, k_2, k_3 \geq 0}.$$
where the $k$-th box from the right carries the parameter $q^{m+1-2k}$. This can be rewritten further as

$$\prod_{j=1}^{m} Y_{1,q^{m+1-2j}z} \sum_{k_1+k_2+k_3=m, k_1,k_2,k_3 \geq 0} k_2 + k_3 \prod_{j=1}^{k_2+k_3} A^{-1}_{1,q^{m+2-2j}z} \prod_{j=1}^{k_3} A^{-1}_{2,q^{m+3-2j}z}$$

$$= \sum_{k_1+k_2+k_3=m, k_1,k_2,k_3 \geq 0} \varepsilon(k_1-k_2)\omega_1 \varepsilon(k_2-k_3)\omega_2 \times \left[ \left( \frac{(1-q^{-m}z)(1-q^{-2m+2}z)}{(1-q^{-2m+2z})(1-q^{-2m+1}z)(1-q^{-m+3}z)} \right) \left( \frac{(1-q^{-m+1}(k_3+k_2)+1)z}{(1-q^{-m+2(k_3+k_2)+1}z)} \right) \left( \frac{(1-q^{-m+2(k_3+k_2)+1}z)(1-q^{-m+3}z)}{(1-q^{-m+2(k_3+k_2)+1}z)} \right) \right].$$

Let us consider the limit $m \to \infty$. There are three possibilities to obtain meaningful answers,

$$k_1 \to \infty, \quad k_2, k_3 : \text{finite}, \quad q^{-k_1} \to 0,$$
$$k_2 \to \infty, \quad k_1, k_3 : \text{finite}, \quad q^{-k_2} \to 0,$$
$$k_3 \to \infty, \quad k_1, k_2 : \text{finite}, \quad q^{-k_3} \to 0.$$

Writing $e^{\alpha_i} = x_i/x_{i+1}$ and defining for $i = 1, 2, 3$

$$Q_{i,z}^{(1)} = \prod_{j=1}^{i-1} (1-x_j/x_i) \times \lim_{q^{-k_i} \to 0} x_i^{-m}(x_1x_2x_3)^{m} \chi_q \left( W_{m,q^{-m}z}^{(1)} \right),$$

we obtain the result

$$Q_{1,z}^{(1)} = Y_{1,z} \sum_{k_2,k_3 \geq 0} \prod_{j=1}^{k_2+k_3} A^{-1}_{1,q^{1-2j}z} \prod_{j=1}^{k_3} A^{-1}_{2,q^{1-2j}z},$$
$$Q_{2,z}^{(1)} = Y_{1,q^2z}^{-1} \sum_{k_3 \geq 0} \prod_{j=1}^{k_3} A^{-1}_{2,q^{3-2j}z}, \quad Q_{3,z}^{(1)} = Y_{2,q^3z}^{-1}.$$

Up to simple multipliers, they are the irreducible $q$-characters

$$\chi_q(L_{1,z}^{-}) = Q_{1,z}^{(1)}, \quad \chi_q(L(Y_{1,q^2z}^{-1}Y_{2,q^2z})) = \frac{Q_{2,z}^{(1)}}{1-e^{-\alpha_1}},$$
$$\chi_q(L_{2,q^3z}^{+}) = \frac{Q_{3,z}^{(1)}}{(1-e^{-\alpha_1})(1-e^{-\alpha_2})}.$$

Using the explicit construction of modules [23], it can be checked that the second one is the $q$-character of the simple module whose highest $\ell$-weight corresponds to the monomial $Y_{1,q^2z}^{-1}Y_{2,q^2z}$.

We give two more examples.

**Example 4.2** ($\mathfrak{g} = B_2^{(1)}$).

$$Q_{1,z}^{(1)} = Y_{1,z} \left\{ \sum_{k_2,k_3,k_1 \geq 0} \prod_{l=1}^{K_1} A^{-1}_{1,q^{1-2l}z} \prod_{l=1}^{K_2} A^{-1}_{2,q^{2-2l}z} \prod_{l=1}^{K_2} A^{-1}_{2,q^{2-2l}z} \prod_{l=1}^{K_2} A^{-1}_{2,q^{2-2l}z} \prod_{l=1}^{K_2} A^{-1}_{2,q^{2-2l}z} \right\}.$$
Example 4.3 (g = C\(_2\)\(^{(1)}\)).

\[ Q_{2,z}^{(1)} = Y_{1,q^2}^{-1} Y_{2,q^2} \left\{ \sum_{k_2,k_1 \geq 0} \prod_{l=1}^{K_1} A_{1,q^3-2l}^{-1} \prod_{l=1}^{K_2} A_{2,q^2-2l}^{-1} \prod_{l=1}^{K_3} A_{2,q^3-2l}^{-1} \right\} , \]

\[ Q_{2,z}^{(1)} = Y_{1,q^2} Y_{2,q^2} \sum_{k_2,k_1 \geq 0} \prod_{l=1}^{K_1} A_{1,q^3-2l}^{-1} \prod_{l=1}^{K_2} A_{2,q^2-2l}^{-1} \prod_{l=1}^{K_3} A_{2,q^3-2l}^{-1} , \]

\[ Q_{0,z}^{(1)} = \tilde{Q}_{0,z}^{(1)} + \tilde{Q}_{0,z}^{(1)} , \]

\[ Q_{0,z}^{(1)} = \tilde{Q}_{0,z}^{(1)} , \]

\[ Q_{0,z}^{(1)} = x_2^{-1} Y_{1,q^2} Y_{1,q^2} Y_{2,q^2} Y_{2,q^2} \sum_{k_2,k_1 \geq 0, K_3 \text{ even}} \prod_{l=1}^{K_1} A_{1,q^3-2l}^{-1} \prod_{l=1}^{K_2} A_{2,q^2-2l}^{-1} , \]

\[ Q_{2,z}^{(1)} = Y_{1,q^2} Y_{2,q^2} \sum_{K_1 \text{ odd}} \prod_{l=1}^{K_1} A_{1,q^3-2l}^{-1} \prod_{l=1}^{K_2} A_{2,q^2-2l}^{-1} , \]

\[ Q_{1,z}^{(1)} = Y_{1,q^3} . \]

Here we set \( K_1 = k_1, K_2 = k_2 + k_1 \) and \( K_2 = k_2 + k_2 + k_1 \). We have

\[ \chi_q(L_{1,z}^-) = Q_{1,z}^{(1)} \quad \text{ and } \quad \chi_q(L_{1,q^3}^+) = \frac{Q_{1,z}^{(1)}}{(1-e^{-\alpha_1})(1-e^{-\alpha_1-\alpha_2})(1-e^{-2\alpha_1-2\alpha_2})} . \]

4.2 Series \( Q_{i,z}^{(1)} \)

In order to discuss the general case, let us prepare some notation. For \( g = X_n^{(1)} (X = A, B, C, D) \), introduce a parametrization of \( \{\alpha_i, \omega_i\} \) by orthonormal vectors \( \{\epsilon_i\} \), and an index set \( J \) with
a partial ordering $\prec$ as follows.

**$A_n^{(1)}$:**

\[
\alpha_i = \epsilon_i - \epsilon_{i+1} \quad (1 \leq i \leq n),
\]

\[
\omega_i = \epsilon_1 + \cdots + \epsilon_i - \frac{i}{n+1} \sum_{j=1}^{n+1} \epsilon_j \quad (1 \leq i \leq n),
\]

\[J = \{1, 2, \ldots, n+1\}, \quad 1 \prec 2 \prec \cdots \prec n+1,\]

**$B_n^{(1)}$:**

\[
\alpha_i = \epsilon_i - \epsilon_{i+1} \quad (1 \leq i \leq n-1), \quad \alpha_n = \epsilon_n,
\]

\[
\omega_i = \epsilon_1 + \cdots + \epsilon_i \quad (1 \leq i \leq n-1), \quad \omega_n = \frac{1}{2}(\epsilon_1 + \cdots + \epsilon_n),
\]

\[J = \{1, \ldots, n, \bar{n}, \ldots, \bar{1}\}, \quad 1 \prec \cdots \prec n < \bar{n} \prec \cdots \prec \bar{1},\]

**$C_n^{(1)}$:**

\[
\alpha_i = \frac{1}{\sqrt{2}}(\epsilon_i - \epsilon_{i+1}) \quad (1 \leq i \leq n-1), \quad \alpha_n = \sqrt{2} \epsilon_n,
\]

\[
\omega_i = \frac{1}{\sqrt{2}}(\epsilon_1 + \cdots + \epsilon_i) \quad (1 \leq i \leq n),
\]

\[J = \{1, \ldots, n, 0, \bar{n}, \ldots, \bar{1}\}, \quad 1 \prec \cdots \prec n < 0 < \bar{n} \prec \cdots \prec \bar{1},\]

**$D_n^{(1)}$:**

\[
\alpha_i = \epsilon_i - \epsilon_{i+1} \quad (1 \leq i \leq n-1), \quad \alpha_n = \epsilon_{n-1} + \epsilon_n,
\]

\[
\omega_i = \epsilon_1 + \cdots + \epsilon_i \quad (1 \leq i \leq n-2), \quad \omega_{n-1} = \frac{1}{2}(\epsilon_1 + \cdots + \epsilon_{n-1} - \epsilon_n), \quad \omega_n = \frac{1}{2}(\epsilon_1 + \cdots + \epsilon_{n-1} + \epsilon_n),
\]

\[J = \{1, \ldots, n, \bar{n}, \ldots, \bar{1}\}, \quad 1 \prec \cdots \prec n-1 < \frac{\bar{n}}{n} < \frac{n-1}{\bar{n}} \prec \cdots \prec \bar{1}.\]

Define also $x_i$ and $f_{j,k}$ for $j, k \in J$ by

**$A_n^{(1)}$:**

\[
x_i = e^{\epsilon_i} \quad (1 \leq i \leq n+1), \quad f_{j,k} = \frac{1}{1-x_k/x_j},
\]

**$B_n^{(1)}$:**

\[
x_i = e^{\epsilon_i} = x_i^{-1} \quad (1 \leq i \leq n), \quad f_{j,k} = \frac{1 + \delta_{k,j}}{1 - x_k/x_j},
\]

**$C_n^{(1)}$:**

\[
x_i = e^{\sqrt{2} \epsilon_i} = x_i^{-1} \quad (1 \leq i \leq n), \quad x_0 = 1,
\]

\[f_{j,k} = \frac{1}{(1-x_k/x_j)(1 + \delta_{k,0}x_k/x_j)},\]

**$D_n^{(1)}$:**

\[
x_i = e^{\epsilon_i} = x_i^{-1} \quad (1 \leq i \leq n), \quad f_{j,k} = \frac{1 - \delta_{k,j}x_k/x_j}{1 - x_k/x_j}.
\]

In the following, in the sum of the form $\sum_{k_i, \ldots, k_j}$, unless mentioned explicitly, $k_i, \ldots, k_j$ run over all non-negative integers, and we use the abbreviation $K_l = \sum_{\mu=l}^{j} k_{\mu}$ for $i < l < j$. In the case $g = C_n^{(1)}$ and $l < 0$, we set $K_l = \sum_{j=l}^{n} k_j + 2k_0 + \sum_{j=1}^{n} k_j$.

We give below the formula for $Q_{i,z}^{(1)}$ for each $i \in J$.

**Case $A_n^{(1)}$:** For $1 \leq i \leq n+1$,

\[
Q_{i,z}^{(1)} = \Phi_{i,z} \sum_{k_{i+1}, \ldots, k_{n+1}} \prod_{j=i}^{n} K_{j+1} A_{j,q_{j+1}i-2z}^{-1}.
\]

where

\[
\Phi_{i,z} = \gamma_{i-1,q_{i-1}z}^{-1} \gamma_{i,q_{i-1}z}^{-1}. \quad (4.4)
\]
Case $D_n^{(1)}$: For $1 \leq i \leq n$,

$$Q_{i,z}^{(1)} = \Phi_{i,z} \left\{ \sum_{k_{i+1}, \ldots, k_n, k_0, k_1} \prod_{k_{j+1}, \ldots, k_{j+1}} A_{j,q^{i+1-2l_z}}^{-1} \prod_{j=1}^{n} A_{j,q^{2n-j-2l_z}}^{-1} \right\},$$

where

$$
\Phi_{i,z} = \mathcal{Y}_{i-1,q^{i+z}}^{-1} \mathcal{Y}_{i,q^{i}}^{-1}, \\
\Phi_{i,z} = \mathcal{Y}_{i-1,q^{2n-i-1}} \mathcal{Y}_{i,q^{2n-i}}^{-1}.
$$

Case $E_n^{(1)}$:

$$Q_{i,z}^{(1)} = \Phi_{i,z} \sum_{k_{i+1}, \ldots, k_n, k_0, k_1} \prod_{k_{j+1}, \ldots, k_{j+1}} A_{j,q^{i+1-2l_z}}^{-1} \prod_{j=1}^{n} A_{j,q^{2n-j-2l_z}}^{-1},$$

where

$$
\Phi_{i,z} = \mathcal{Y}_{i-1,q^{i+z}}^{-1} \mathcal{Y}_{i,q^{i}}^{-1}, \\
\Phi_{i,z} = \mathcal{Y}_{i-1,q^{2n-i-1}} \mathcal{Y}_{i,q^{2n-i}}^{-1}.
$$
In this subsection we give the definition of \( Q \). We set
\[
\begin{align*}
J, n, \emptyset (0) & = 1, z (1) = J, z (1), \\
n, z & = A^{-1}_{j, q^{2n-1} - j - 2l_z} A^{-1}_{n, q^{n-2l_z}} A^{-1}_{j, q^{j+1} - 1 - 2l_z} \\
1 & \leq i \leq n - 1,
\end{align*}
\]
where
\[
\begin{align*}
Q_{n, z}^{(1)} & = \Phi_{n, z} \sum_{k_{n-1}, \ldots, k_1} K_j A^{-1}_{j, q^{2n-1} - j - 2l_z} A^{-1}_{n, q^{n-2l_z}} A^{-1}_{j, q^{j+1} - 1 - 2l_z}, \\
Q_{l, z}^{(1)} & = \Phi_{l, z} \sum_{k_{n-1}, \ldots, k_1} K_j A^{-1}_{j, q^{2n-1} - j - 2l_z}, 1 \leq i \leq n,
\end{align*}
\]

**Remark 4.4.** As in [26, 27], it is known that the \( q \)-character of KR module \( W_{m, z}^{(1)} \) for \( D_{n}^{(1)} \) can be written in terms of tableau of elementary boxes \([a], a \in \{1, \ldots, n, 0, \tilde{n}, \ldots, \tilde{1}\}\), while the elementary box \([\emptyset]\) is allowed to appear at most once. When taking limit of the \( q \)-character of \( W_{m, z}^{(1)} \), we have only \( 2n \) kinds of meaningful results, then we use the index set \( \{1, \ldots, n, \tilde{n}, \ldots, \tilde{1}\}\). By similar reason we use \( \{1, \ldots, n, 0, \tilde{n}, \ldots, \tilde{1}\}\) as the index set of \( C_{n}^{(1)} \), which looks like opposite to the usual one.

### 4.3 Series \( Q_{J, z}^{(a)} \)

In this subsection we give the definition of \( Q_{J, z}^{(a)} \) where \( J \in J^a \) and \( (g, a) \neq (C_{n}^{(1)}, n), (D_{n}^{(1)}, n-1), (D_{n}^{(1)}, n) \). We recall that \( q_1 = q^{1/2} \) for \( C_{n}^{(1)} \) and \( q_1 = q \) in the other cases.

We fix our convention about the indices as follows. For \( J = (j_1, \ldots, j_a) \in J^a \), we denote by \( J_\sim \) the underlying set \( \{j_1, \ldots, j_a\} \subseteq J \). Set further
\[
J = (\overline{j_1}, \ldots, \overline{j_1}), \\
J_\sim = (i_1, \ldots, i_b), \quad \text{with } J_\sim = J \setminus J_\sim \quad \text{and} \quad i_1 < \cdots < i_b.
\]

We say \( J \) is increasing if \( j_1 < \cdots < j_a \).

For an element \( J = (j_1, \ldots, j_a) \in J^a \), we define \( Q_{J, z}^{(a)} \) by
\[
Q_{J, z}^{(a)} = \det \left( Q_{j_\mu, q_1^{j_\nu} x_\nu}^{(1)} \right)_{\mu, \nu = 1}^{a}
\]
\[
\begin{pmatrix}
Q_{j_1, q_1^{j_1} x_1}^{(1)} & \cdots & Q_{j_a, q_1^{j_a} x_a}^{(1)} \\
Q_{j_1, q_1^{j_1} x_1}^{(1)} & \cdots & Q_{j_a, q_1^{j_a} x_a}^{(1)} \\
\vdots & \ddots & \vdots \\
Q_{j_1, q_1^{j_1} x_1}^{(1)} & \cdots & Q_{j_a, q_1^{j_a} x_a}^{(1)} \\
Q_{j_1, q_1^{j_1} x_1}^{(1)} & \cdots & Q_{j_a, q_1^{j_a} x_a}^{(1)} \\
Q_{j_1, q_1^{j_1} x_1}^{(1)} & \cdots & Q_{j_a, q_1^{j_a} x_a}^{(1)} \\
\end{pmatrix}
\]

We set \( Q_{\emptyset, z}^{(0)} = 1 \).
Note that if the entries of $J \in \mathcal{J}^a$ are not distinct (i.e. $\forall j < a$), then $Q_{J,z}^{(a)} = 0$.

Define
\[ \tilde{Q}_{J}^{(a)} = \varpi(Q_{J,z}^{(a)}). \]

Hence by the definition of $Q_{J,z}^{(a)}$, it is easy to get
\[ \tilde{Q}_{J}^{(a)} = \prod_{i \in J, j \in J^*} f_{i,j} \prod_{i \in J, i < j} ((x_j - x_i)f_{i,j}). \tag{4.15} \]

For an increasing element $J = (j_1, \ldots, j_a) \in \mathcal{J}^a$, set
\[ \Phi_{J,z} = \prod_{k=1}^{a} \Phi_{j_k, q_k^{a+1-2k} z}, \tag{4.16} \]

where $\Phi_{J,z}$ are defined in (4.4)–(4.11) with $Y_{i,z}$ being given in (4.1). Let $L(\Phi_{J,z})$ be the unique irreducible $U_q(b)$-module with the highest $\ell$-weight $\Phi_{J,z}$. Except in the case $\mathfrak{g} = C_{n}^{(1)}$, we expect that the $q$-character of $L(\Phi_{J,z})$ is given by
\[ \chi_q(L(\Phi_{J,z})) = \prod_{i \in J, j \in J^*} f_{j,i} \prod_{i \in J, i < j} (-x_i^{-1}) \cdot Q_{J,z}^{(a)}. \]

In particular, from (4.15) we expect that
\[ \chi(L(\Phi_{J,z})) = \prod_{(i,j) \in (J \times J^*) \cup (J^* \times J), i \prec j} f_{i,j} \times \begin{cases} 1, & \mathfrak{g} = A_{n}^{(1)}, \\ \prod_{i \in J, i \prec j} (1 + x_i^{-1}), & \mathfrak{g} = B_{n}^{(1)}, \\ \prod_{i \in J, i \prec j} (1 - x_i^{-2}), & \mathfrak{g} = D_{n}^{(1)}. \end{cases} \]

In the case of $\mathfrak{g} = C_{n}^{(1)}$, we expect the same to be true for $Q_{J,z}^{(a)}$ if $j_1, \ldots, j_a \neq 0$. As noted already in Example 4.3, the series $Q_{0,z}^{(1)}$ does not correspond to an irreducible $q$-character. We expect rather that $\tilde{Q}_{0,z}^{(1)}$, $\tilde{Q}_{0,z}^{(1)}$ correspond to irreducible $q$-characters. See also Remark 5.9.

### 4.4 Series $\mathcal{R}_{\varepsilon,z}^{(n)}$ for the spin node: case $C_{n}^{(1)}$

In this subsection, we introduce another set of series $\mathcal{R}_{\varepsilon,z}^{(n)}$ for the spin node. Unlike the series $Q_{J,z}^{(a)}$, in general we do not know the explicit formulas for them. We define $\mathcal{R}_{\varepsilon,z}^{(n)}$ by the $q$-characters of the irreducible $U_q(b)$ module $L(M_{\varepsilon,z})$, and give a rule to determine the highest $\ell$-weight $M_{\varepsilon,z}$ which is a monomial in $\mathbb{C}[y_{i,z}^{\pm 1}]_{1 \leq i, z \in \mathbb{C}^*}$.

Exhibiting the $n$-dependence explicitly, let $P_n$ be the weight lattice of the simple Lie algebra of type $C_n$. Set
\[ E_n = \{ \varepsilon = (\varepsilon_1, \ldots, \varepsilon_n) \mid \varepsilon_i = \pm 1, 1 \leq i \leq n \}. \]

We define two weight functions
\[ w_1 : E_n \rightarrow P_n, \quad \varepsilon \mapsto \frac{1}{2} \sum_{i=1}^{n} \varepsilon_i \epsilon_i, \tag{4.17} \]
\[ w_2 : \{ J \mid J \in \mathcal{J}^a(1 \leq a \leq \#\mathcal{J}), J \subseteq \mathcal{J} \} \rightarrow P_n, \quad (j_1, \ldots, j_n) \mapsto \sum_{k=1}^{a} \text{sgn}(j_k)\epsilon_{j_k}, \quad (4.18) \]

where \( \epsilon_0 = 0 \), and for \( j \in \mathcal{J} \) we set

\[ \text{sgn}(j) = \begin{cases} +, & j \leq n, \\ 0, & j = 0, \\ -, & j \geq n. \end{cases} \]

Borrowing an idea from [26], we introduce monomials \( M_{\epsilon,z} \) in \( \mathbb{C}[\mathcal{Y}_{i,z}]_{1 \leq i \leq n,z \in \mathbb{C}^\times} \) inductively as follows. Define two operators \( \tau^\mathcal{Y}, \tau^\mathcal{Z}_c \) by

\[ \tau^\mathcal{Y} : \mathcal{Y}_{i,z} \rightarrow \mathcal{Y}_{i+1,z}, \quad \tau^\mathcal{Z}_c : \mathcal{Y}_{i,z} \rightarrow \mathcal{Y}_{1,q^2z}. \]

For \( n = 1 \) we set \( M_{+,z} = \mathcal{Y}_{1,z} \) and \( M_{-,z} = \mathcal{Y}_{1,q^2z}^{-1} \). In the general case we set

\[ \begin{align*}
M_{(+,-),z} &= \tau^\mathcal{Y}(M_{(+,-),z}), \\
M_{(-,-),z} &= \mathcal{Y}_{1,q^2z}^{-1} \tau^\mathcal{Y}(M_{(-,-),z}), \\
M_{(+,+),z} &= \mathcal{Y}_{1,q^2z} \tau^\mathcal{Y}(M_{(+,+),z}), \\
M_{(-,+),z} &= \tau^\mathcal{Z}_c \tau^\mathcal{Y}(M_{(-,+),z}),
\end{align*} \]

(4.19)

where \( \xi \in E_{n-2} \).

Now let \( \epsilon \in E_n \) with \( w_1(\epsilon) = \frac{1}{2} (\sum_{k=1}^{n} \epsilon \varphi_k - \sum_{k=1}^{n} \epsilon \psi_k) \), and consider the simple module \( L(M_{\epsilon,z}) \). We define \( \mathcal{R}_{\epsilon,z}^{(n)} \) and \( \mathcal{R}_{\epsilon}^{(n)} \) through the \( q \)-character and the character as follows.

\[ \begin{align*}
\mathcal{R}_{\epsilon,z}^{(n)} &= \chi_q(L(M_{\epsilon,z})) \prod_{1 \leq k \leq l \leq t} (1 - x_{\psi_k}x_{\psi_l}) \prod_{1 \leq k < s \leq l \leq t, \varphi_k > \varphi_l}(1 - x_{\varphi_k}^{-1}x_{\psi_l}), \\
\mathcal{R}_{\epsilon}^{(n)} &= \chi(L(M_{\epsilon,z})) \prod_{1 \leq k \leq l \leq t} (1 - x_{\psi_k}x_{\psi_l}) \prod_{1 \leq k < s \leq l \leq t, \varphi_k > \varphi_l}(1 - x_{\varphi_k}^{-1}x_{\psi_l}).
\end{align*} \]

For the latter we have the following guess.

**Conjecture 4.5.** For \( \epsilon = (\epsilon_1, \ldots, \epsilon_n) \in E_n \), we have

\[ \chi(L(M_{\epsilon,z})) = \prod_{1 \leq i \leq n} \frac{1}{1 - x_i^{-2\epsilon_i}} \prod_{1 \leq i < j \leq n} \frac{1}{1 - x_i^{-\epsilon_i}x_j^{-\epsilon_j}}. \]

In the special case \( n = 2 \), one can obtain the series for \( C_{2}^{(1)} \) from that of \( B_{2}^{(1)} \). Let \( Q_{1,z}^{(1)}(1 \leftrightarrow 2) \) stand for the series for \( B_{2}^{(1)} \) with \( a = 1 \) given in Example 4.2, wherein we interchange \( \mathcal{Y}_{1,z} \) with \( \mathcal{Y}_{2,z} \) and \( A_{1,z} \) with \( A_{2,z} \). Then we have

\[ \begin{align*}
\mathcal{R}_{(+,+),z}^{(2)} &= Q_{1,z}^{(1)}(1 \leftrightarrow 2), \\
\mathcal{R}_{(+,-),z}^{(2)} &= Q_{2,z}^{(1)}(1 \leftrightarrow 2), \\
\mathcal{R}_{(-,+),z}^{(2)} &= Q_{2,z}^{(1)}(1 \leftrightarrow 2), \\
\mathcal{R}_{(-,-),z}^{(2)} &= Q_{1,z}^{(1)}(1 \leftrightarrow 2).
\end{align*} \]

(4.21)

**4.5 Series \( \mathcal{R}_{\epsilon,z}^{(n-1)}, \mathcal{R}_{\epsilon}^{(n)} \) for the spin nodes: case \( D_{n}^{(1)} \)**

For the two spin nodes of \( D_{n}^{(1)} \), we follow the procedure for \( C_{n}^{(1)} \). We use the same weight functions (4.17), (4.18) as for \( C_{n}^{(1)} \). We also introduce two subsets of \( E_n \),

\[ E_{n,c} = \{ \epsilon \in E_n \mid x^c \{ \epsilon_i = -1 \} = \varsigma \pmod{2} \}, \quad \varsigma = 0, 1. \]
We define monomials $M_{\varepsilon,z}$ inductively by
\[ M_{(+\varepsilon),z} = \tau^Y(M_{+\varepsilon},z), \quad M_{(-\varepsilon),z} = \mathcal{Y}_{1,q^2}^z \tau^Y(M_{-\varepsilon},z), \] (4.23)
\[ M_{(\varepsilon),z} = \mathcal{Y}_{1,q}^z \tau^Y(M_{+\varepsilon},z), \quad M_{(-\varepsilon),z} = \tau^Y(M_{-\varepsilon},z), \] (4.24)
with $\varepsilon \in E_{n-2}$ and the initial values
\[ M_{(+\varepsilon),z} = \mathcal{Y}_{2,z}, \quad M_{(-\varepsilon),z} = \mathcal{Y}_{2,q^2}^z, \]
\[ M_{(\varepsilon),z} = \mathcal{X}_{1,z}, \quad M_{(-\varepsilon),z} = \mathcal{Y}_{1,q}^{-1}z. \]

For $\varepsilon \in E_{n,\zeta}$ ($\zeta = 0, 1$) with $w_1(\varepsilon) = \frac{1}{2} \left( \sum_{k=1}^s \varepsilon_{\psi_k} - \sum_{l=1}^t \varepsilon_{\psi_l} \right)$, we define
\[ \mathcal{R}_{\varepsilon,z}^{(n-\varepsilon)} = \mathcal{X}_q(L(M_{\varepsilon,z})) \prod_{1 \leq k < l \leq t} (1 - x_{\psi_k} x_{\psi_l}) \prod_{1 \leq k < l \leq t, \psi_k > \psi_l} (1 - x_{\psi_k}^{-1} x_{\psi_l}), \]
\[ \mathcal{R}_{\varepsilon,z}^{(-n-\varepsilon)} = \mathcal{X}_q(L(M_{\varepsilon,z})) \prod_{1 \leq k < l \leq t} (1 - x_{\psi_k} x_{\psi_l}) \prod_{1 \leq k < l \leq t, \psi_k > \psi_l} (1 - x_{\psi_k}^{-1} x_{\psi_l}). \]

**Conjecture 4.6.** For $\varepsilon = (\varepsilon_1, \ldots, \varepsilon_n) \in E_{n,\zeta}$ ($\zeta = 0, 1$), we have
\[ \mathcal{X}_q(L(M_{\varepsilon,z})) = \prod_{1 \leq i < j \leq n} 1 - x_i^{-\varepsilon_i} x_j^{-\varepsilon_j}. \]

Although we do not have formulas for the series $\mathcal{R}_{\varepsilon,z}^{(n-\varepsilon)}$ ($\zeta = 0, 1$) for general $n$, in the special case $D_k^{(1)}$, similarly to $Q_{i,z}^{(1)}$, we define $\mathcal{R}_{\varepsilon,z}^{(3)}$ (resp. $\mathcal{R}_{\varepsilon,z}^{(4)}$) as certain limits of the q-characters of the KR modules $W_{m,z}^{(3)}$ (resp. $W_{m,z}^{(4)}$). More precisely, one can obtain $\mathcal{R}_{\varepsilon,z}^{(3)}$ (resp. $\mathcal{R}_{\varepsilon,z}^{(4)}$) from $Q_{i,z}^{(1)}$ by exchanging $\mathcal{X}_{1,z}$ with $\mathcal{X}_{3,z}$ (resp. $\mathcal{X}_{4,z}$) and $A_{1,z}$ with $A_{3,z}$ (resp. $A_{4,z}$). For example, we have
\[ Q_{i,z}^{(1)} = \mathcal{Y}_{2,q^2}^{-1} \sum_{k=0}^k A_{1,q^6-2j-z}^{-1} \prod_{j=1}^k, \quad Q_{i,z}^{(1)} = \mathcal{Y}_{1,q^6}^{-1}, \]
\[ R_{(-++)},z = \mathcal{Y}_{2,q^2}^{-1} \sum_{k=0}^k A_{3,q^6-2j-z}^{-1}, \quad R_{(-++)},z = \mathcal{Y}_{3,q^6}^{-1}, \]
\[ R_{(++)},z = \mathcal{Y}_{2,q^2}^{-1} \sum_{k=0}^k A_{4,q^6-2j-z}^{-1}, \quad R_{(++)},z = \mathcal{Y}_{4,q^6}^{-1}, \]
and so on.

## 5 Polynomial relations

In this section, we shall give the main results of this paper. We propose polynomial relations among the series defined in Section 4, which are expected from the $\psi$-system obtained in Section 3.

We prepare some notation which we will use below. For an element $j \in \mathcal{J}$, an element $J = (j_1, \ldots, j_a) \in \mathcal{J}^a$ and a positive integer $k$ with $1 \leq k \leq a$, we set
\[ \lambda_j = \sum_{i=1}^a \lambda_{j_i}, \quad x_J = \prod_{i=1}^a x_{j_i}, \quad J - j_k = (j_1, \ldots, \hat{j_k}, \ldots, j_a), \quad (j, J) = (j, j_1, \ldots, j_a). \]

Let $\{v_i\}_{i \in \bar{\mathcal{J}}}$ be the standard basis of the vector representation $V^{(1)}$ of $L_{\mathfrak{g}}$ (cf. Appendix A). Here $\bar{\mathcal{J}} = \{1, \ldots, n, 0, \bar{n}, \ldots, \bar{1}\}$ for $C_n^{(1)}$ and $\bar{\mathcal{J}} = \mathcal{J}$ for the other cases. For $J = (j_1, \ldots, j_a) \in \bar{\mathcal{J}}^a$, we set $v_J = v_{j_1} \wedge v_{j_2} \wedge \cdots \wedge v_{j_a}$. 

5.1 Plücker-type relations

We begin by writing down the \( \psi \)-system and the corresponding relations for \( Q_{J,z}^{(a)} \) for all nodes excepting the last (or the second last for \( D_n^{(1)} \)). In fact they are just the consequence of the fact that the \( Q_{J,z}^{(a)} \) are defined by determinants of the \( Q_{i,z}^{(1)} \)'s. So they are rather definitions, and we write them down just for uniformity reasons. As we shall see later, the only non-trivial relations come from the last (or the second last) node.

For comparison, we start with the \( \psi \)-system. Let \( J_1 = (i_1, \ldots, i_a) \), \( J_2 = (j_1, \ldots, j_a) \) be two elements of \( J^a \). The embedding \( \iota (3.11) \) of \( L_\mathfrak{g} \)-modules is explicitly given by

\[
\iota : \bigwedge^2 V_2^{(a)} \hookrightarrow V^{(a-1)} \otimes V^{(a+1)},
\]

\[
v_{J_1} \wedge v_{J_2} \mapsto \frac{1}{2} \left( \sum_{k=1}^{a} (-1)^{k-1} v_{J_1} \otimes v_{(i_k, J_2)} - \sum_{k=1}^{a} (-1)^{k-1} v_{J_2} \otimes v_{(j_k, J_1)} \right).
\]

Therefore (3.12), (3.14), (3.16) and (3.19) imply

\[
Q_{J_1, J_2}^{(a)} (E) = Q_{J_1, J_2, J_2}^{(a)} (E) - Q_{J_1, J_2, J_2}^{(a)} (E) \omega^{\lambda_j - \lambda_j^2} = \sum_{k=1}^{a} (-1)^k Q_{J_2, J_2}^{(a-1)} (E) Q_{(j_k, J_1)}^{(a+1)} (E) \omega^{-\lambda_{jk}}.
\]

The corresponding relations for \( Q_{J,z}^{(a)} \) read as follows.

**Proposition 5.1.** For two increasing elements \( J_1 = (i_1, i_2, \ldots, i_a) \) and \( J_2 = (j_1, j_2, \ldots, j_a) \) of \( J^a \), we have

\[
Q_{J_1, q_1, J_2, q_2, z}^{(a)} - Q_{J_1, q_1, J_2, q_2, z}^{(a)} = \sum_{k=1}^{a} (-1)^k Q_{J_2, J_2}^{(a-1)} z Q_{(j_k, J_1)}^{(a+1)} z^{-1} z^{-1}.
\]

This is an immediate consequence of the Sylvester identity (see e.g. [18, p. 108]). Note that (5.1) and (5.2) has the same form under the identification

\[
Q^{(a)}_J (E) \leftrightarrow Q^{(a)}_{J,z}, \quad E \leftrightarrow z, \quad \Omega^{1/2} \leftrightarrow q_1, \quad \omega \leftrightarrow \times_i.
\]

5.2 Wronskian identity for \( A_n^{(1)} \)

For type \( A_n^{(1)} \), the non-trivial relation is the following ‘Wronskian identity’.

**Theorem 5.2.** Let \( \mathfrak{g} = A_n^{(1)} \). We have

\[
\det \left( Q_{\nu, q}^{(1)} n^{n+2-2\mu, x_{\mu}, \mu, \nu} \right)_{\mu, \nu = 1}^{n+1} = 1.
\]

Note that the left-hand side is just \( Q_{J,z}^{(n+1)} (-1)^{(n+1)/2} \prod_{b=1}^{n} x_b^{n+1-b} \) with \( J = (1, 2, \ldots, n+1) \). We prove Theorem 5.2 in Appendix B.
5.3 Polynomial relations related to the last node: case $B_n^{(1)}$

In this subsection we let $g = B_n^{(1)}$. We give conjectural relations which correspond to the last identity (3.15) of the $\psi$-system. We first give identities for the connection coefficients $Q_{J}^{(a)}(E)$.

In this subsection, for an element $j \in J$ and $J = (j_1, \ldots, j_a) \in J^n$ we define

$$c(j) = \begin{cases} (-1)^j, & j \leq n, \\ (-1)^{j-1}, & j > n, \end{cases} \quad c(J) = \prod_{k=1}^{a} c(j_k).$$

We also denote by $\sigma(J, J^*)$ the signature of the permutation $(1, \ldots, n, \bar{n}, \ldots, \bar{1}) \mapsto (J, J^*)$.

The bilinear form $\langle \cdot, \cdot \rangle: V \times V \to \mathbb{C}$ given by $\langle v_i, v_j \rangle = c(i)\delta_{i,j}$ is $L_g$-invariant. From this it is easy to see that we have the following $L_g$-module isomorphism

$$V_{1/2}^{(a)} \simeq V^{(2n-a)}, \quad a = 1, \ldots, n, \quad v_J \mapsto c(J)\sigma(J, J^*)v_{J^*},$$

which implies

$$Q_{J,\frac{1}{2}}^{(a)}(E) = c(J)\sigma(J, J^*)Q_{T^*}^{(2n-a)}(E). \quad (5.3)$$

Let $J_1 = (i_1, \ldots, i_n), J_2 = (j_1, \ldots, j_n)$ be two elements of $J^n$. Then the relation (5.1) together with (5.3) imply

$$Q_{J_1,\frac{1}{2}}^{(n)}(E)Q_{J_2,\frac{1}{2}}^{(n)}(E)c \left( \frac{J_1}{J_2} \right) \sigma \left( \frac{J_1}{J_2}, J_1 \right) \omega^{-\lambda_{J_1}} = Q_{J_2,\frac{1}{2}}^{(n)}(E)Q_{J_1,\frac{1}{2}}^{(n)}(E)c \left( \frac{J_1}{J_2}, J_2 \right) \sigma \left( \frac{J_1}{J_2}, J_2 \right) \omega^{-\lambda_{J_2}}$$

$$= \sum_{k=1}^{n} (-1)^k Q_{J_2-j_k}^{(n-1)}(E)Q_{J_1-J_k,\frac{1}{2}}^{(n-1)}(E)c \left( \frac{J_1}{J_2}, J_1 \right) \sigma \left( \frac{J_1}{J_2}, J_1 \right) \omega^{-\lambda_{J_1}-\lambda_{J_k}}. \quad (5.4)$$

For a given element $i \in J$ and $J \in J^n$, let

$$X_J = \prod_{\mu \in J, \mu < i} \left( -x_{\mu}^{1/2}\delta_{\mu, i}^{-1} \right), \quad x_{i, J} = \prod_{j \in J, j < i} \left( -x_j^{1/2}\delta_{j, i}^{-1} \right). \quad (5.5)$$

The counterpart of (5.4) for the series $Q_{J, z}$ is given by the following conjecture.

**Conjecture 5.3.** Let $J_1 = (i_1, \ldots, i_n), J_2 = (j_1, \ldots, j_n) \in J^n$ be increasing. Then we have

$$Q_{J_1, z}^{(n)}Q_{J_2, qz}^{(n)}X_{J_1} X_{J_2} x_{J_1}^{-1/2} - Q_{J_1, qz}^{(n)}Q_{J_2, z}^{(n)}X_{J_1} X_{J_2} x_{J_2}^{-1/2}$$

$$= -\sum_{j_k \in J_{2-j_k}, z \leq k \leq n} Q_{J_2-j_k, qz}^{(n-1)}Q_{J_1, \frac{1}{2}-J_k, qz}^{(n-1)}X_{J_2-j_k} X_{J_1} x_{J_2}^{-1/2} - x_{J_2}^{-1/2} x_{J_1}^{-1/2}. \quad (5.6)$$

When $n = 2$, we have verified the above conjecture by direct computations. To save space we do not write the proofs here. We have also checked them for $n = 3$ by Mathematica 5.0 up to certain degree, counting the degree of $A_{4, q^2z}$ to be $k$. For $n \geq 4$, although at the moment we do not have a proof of these identities, in Appendix C, we show that the identities hold when specialized to the characters.

**Remark 5.4.** It is worth noting that, although the connection coefficients satisfy the linear relation (5.3), there are no analogs for the series $Q_{J, z}^{(a)}$. For example, in the simplest case $B_2^{(1)}$, a direct computation gives $Q_{(221), z}^{(3)} = -x_2 y_1^{-1} y_2 y_1 y_2 y_2 y_1^{-1}, y_2^{-1} y_2 y_1 y_2 y_2 y_1^{-1}; y_1^{-1} y_2, y_1^{-1} y_2 y_2 y_1^{-1}$, and there is no function $f(x_1, x_2)$ of $x_1, x_2$ such that $Q_{(221), z}^{(3)} = f(x_1, x_2)Q_{1, z}^{(1)}$. So only the quadratic relations survive.
5.4 Polynomial relations related to spin node: case $C_n^{(1)}$

In this subsection, $g = C_n^{(1)}$. We shall give conjectural relations for the series $Q_{J,z}^{(n)}$ and $R_{\epsilon,z}^{(n)}$.

For an increasing element $J \in \mathcal{J}^n$ we can write

$$J = \{\alpha_1, \ldots, \alpha_m, \beta_1, \ldots, \beta_r, \sigma_1, \ldots, \sigma_u, \bar{\sigma}_1, \ldots, \bar{\sigma}_v \} \cup (J \cap \{0\}),$$

$$J^* = \{\bar{\alpha}_1, \ldots, \bar{\alpha}_m, \bar{\beta}_1, \ldots, \bar{\beta}_r, \eta_1, \ldots, \eta_v, \bar{\eta}_1, \ldots, \bar{\eta}_v \} \cup (J^* \cap \{0\}),$$

where $\alpha_i$, $\beta_i$, $\sigma_i$, $\eta_i$ are mutually distinct elements of $\{1, \ldots, n\}$. Let $\sigma = \{\sigma_1, \ldots, \sigma_u\}$, $\eta = \{\eta_1, \ldots, \eta_v\}$. Note that

$$v = \begin{cases} u + 2, & 0 \in J \text{ and } J \in \mathcal{J}^{n-1}, \\ u + 1, & 0 \in J^* \text{ and } J \in \mathcal{J}^{n-1}, \text{ or } 0 \in J \text{ and } J \in \mathcal{J}^n, \\ u, & 0 \in J \text{ and } J \in \mathcal{J}^n. \end{cases}$$

Let $J \in \mathcal{J}^n$ be increasing. We introduce the following condition for $\epsilon, \epsilon' \in E_n$:

**Condition $C_J$:**

- $w_1(\epsilon) + w_1(\epsilon') = w_2(J)$,
- if $w_1(\epsilon) - w_1(\epsilon') = \sum_{k=1}^{u} \epsilon_{\gamma k} - \sum_{k=1}^{u} \epsilon_{\gamma k}'$, then $\sigma \subseteq \gamma$.

Here $w_1$, $w_2$ are the weight functions defined by (4.17), (4.18) and $\gamma = \{\gamma_1, \ldots, \gamma_s\}$, $\delta = \{\delta_1, \ldots, \delta_l\}$. It is easy to see $\sigma \cup \eta = \gamma \cup \delta$.

In the following Conjectures 5.5–5.8, $g = C_n^{(1)}$, $\sigma$, $\eta$ are disjoint subsets of $\{1, 2, \ldots, n\}$.

**Conjecture 5.5.** Suppose $J \in \mathcal{J}^{n-1}$ is increasing and $0 \notin J$. Then we have

$$(-1)^{u(u-1)/2} Q_{J,z}^{(n-1)} x_{\sigma}^{-2u} x_{\eta}^{-u} \prod_{i,j \in J \atop i < j} (-x_i^{-1}) \prod_{1 \leq k \leq u, 1 \leq l \leq u+2, \sigma_k > \eta_l} \left( \frac{-x_{\sigma k}}{x_{\eta l}} \right)$$

$$= \sum (-1)^{t(t+1)/2} R_{\epsilon,z}^{(n)} R_{\epsilon',z}^{(n)} x_{\gamma}^{-2} x_{\delta}^{-1-t} \prod_{1 \leq k \leq s, 1 \leq l \leq t, \gamma_k > \delta_l} \left( \frac{-x_{\gamma k}}{x_{\delta l}} \right),$$

where the sum is taken over all $\epsilon, \epsilon' \in E_n$ satisfying Condition $C_J$ where $\gamma$, $\delta$ are determined there.

**Conjecture 5.6.** Suppose $J \in \mathcal{J}^{n-1}$ is increasing and $0 \notin J$. Then we have

$$(-1)^{u(u-1)/2} Q_{J,z}^{(n-1)} x_{\sigma}^{-2u} x_{\eta}^{-u} \prod_{i,j \in J \atop i < j} (-x_i^{-1}) \prod_{1 \leq k \leq u, 1 \leq l \leq u+2, \sigma_k > \eta_l} \left( \frac{-x_{\sigma k}}{x_{\eta l}} \right)$$

$$= \sum (-1)^{t(t+1)/2} R_{\epsilon,z}^{(n)} R_{\epsilon',z}^{(n)} x_{\gamma}^{-2} x_{\delta}^{-1-t} \prod_{1 \leq k \leq s, 1 \leq l \leq t, \gamma_k > \delta_l} \left( \frac{-x_{\gamma k}}{x_{\delta l}} \right),$$

where the sum is taken over all $\epsilon, \epsilon' \in E_n$ satisfying Condition $C_J$ where $\gamma$, $\delta$ are determined there.
Conjecture 5.7. Suppose \( J \in J^n \) is increasing and \( 0 \notin J \). Then we have

\[
(-1)^{u(n-1)/2} Q_{J,z}^{(n)} T^{-2u,n-1-u} \prod_{i,j \in J, i < j} (-x_i^{-1}) \prod_{1 \leq k \leq u, 1 \leq \ell \leq u, \sigma_k > \eta_\ell} \left( -\frac{x_{\sigma_k}}{x_{\eta_\ell}} \right) 
\]

\[
= \sum (-1)^{z(n-1)/2} R_{\varepsilon,q}^{(n)} T \prod_{\varepsilon',q R_{\varepsilon',q}^{(n)} Z} x_1^{-2u} x_1^{-t} \prod_{1 \leq k \leq u, 1 \leq \ell \leq t, \gamma_k > \delta_\ell} \left( -\frac{x_{\gamma_k}}{x_{\delta_\ell}} \right),
\]

where the sum is taken over all \( \varepsilon, \varepsilon' \in E_n \) satisfying Condition \( C_J \) where \( \gamma, \delta \) are determined there.

Conjecture 5.8. Suppose \( J \in J^n \) is increasing and \( 0 \in J \). Then we have

\[
(-1)^{u(n-1)/2} Q_{J,z}^{(n)} T^{-2u,n-1-u} \prod_{i,j \in J, i < j} (-x_i^{-1}) \prod_{1 \leq k \leq u, 1 \leq \ell \leq u+1, \sigma_k > \eta_\ell} \left( -\frac{x_{\sigma_k}}{x_{\eta_\ell}} \right) 
\]

\[
= \sum (-1)^{z(n-1)/2} R_{\varepsilon,q}^{(n)} T \prod_{\varepsilon',q R_{\varepsilon',q}^{(n)} Z} x_1^{-2u} x_1^{-t} \prod_{1 \leq k \leq u, 1 \leq \ell \leq t, \gamma_k > \delta_\ell} \left( -\frac{x_{\gamma_k}}{x_{\delta_\ell}} \right),
\]

where the sum is taken over all \( \varepsilon, \varepsilon' \in E_n \) satisfying Condition \( C_J \) where \( \gamma, \delta \) are determined there.

For \( n = 2 \) one can verify these relations using the relation between \( C_2^{(1)} \) and \( B_2^{(1)} \) (see (4.21), (4.22)). For general \( n \), we show in Appendix D the validity of the identities when specialized to the ordinary characters \( \hat{Q}_J^{(a)} \) and \( \hat{R}_{\varepsilon}^{(n)} \).

Remark 5.9. In the conjectured identities given above, \( \hat{Q}_{0,z}^{(1)} \) and \( \hat{Q}_{0,z}^{(1)} \) enter only through the sum \( \hat{Q}_{0,z}^{(1)} = \hat{Q}_{0,z}^{(1)} + \hat{Q}_{0,z}^{(1)} \). However, since the weight 0 of \( V^{(1)} \) has multiplicity 2, it is more natural to consider \( \hat{Q}_{0,z}^{(1)} \) and \( \hat{Q}_{0,z}^{(1)} \) separately. The \( \psi \)-system also suggests that there are identities involving them separately. This is indeed the case for \( n = 2 \) where we have, for example,

\[
\hat{Q}_{1,q}^{(1)} \hat{Q}_{0,q}^{(1)} - x_1 \hat{Q}_{0,q}^{(1)} \hat{Q}_{1,q}^{(1)} = -x_1 R_{1,q}^{(2)} R_{1,q}^{(2)}(++) q^{-1} z \hat{R}_{1,q}^{(2)}(++) q^{-1} z ,
\]

\[
\hat{Q}_{1,q}^{(1)} \hat{Q}_{0,q}^{(1)} - x_1 \hat{Q}_{0,q}^{(1)} \hat{Q}_{1,q}^{(1)} = -x_1 R_{1,q}^{(2)} R_{1,q}^{(2)}(++) q^{-1} z \hat{R}_{1,q}^{(2)}(++) q^{-1} z .
\]

On the other hand, if we define \( \hat{Q}_{J,z}^{(a)} \) as determinants treating 0 and \( \bar{0} \) independently, then when \( J \) contains both 0 and \( \bar{0} \) we have \( \hat{Q}_{J,z}^{(a)} = 0 \) because \( x_0 = x_0 = 1 \). Hence \( \hat{Q}_{J,z}^{(a)} \) cannot be explained as a \( q \)-character, so we must modify the working hypothesis. At the moment we do not know how to do that.

5.5 Polynomial relations related to spin nodes: case \( D_n^{(1)} \)

Finally we consider the case \( g = D_n^{(1)} \). As in the case of \( C_n^{(1)} \), for an increasing element \( J \in J^a \) we can write

\[
\sim = \{ \alpha_1, \ldots, \alpha_m, \bar{\beta}_1, \ldots, \bar{\beta}_r, \sigma_1, \ldots, \sigma_u, \bar{\sigma}_1, \ldots, \bar{\sigma}_u \},
\]

\[
J^* = \{ \bar{\alpha}_1, \ldots, \bar{\alpha}_m, \bar{\beta}_1, \ldots, \bar{\beta}_r, \eta_1, \ldots, \eta_v, \bar{\eta}_1, \ldots, \bar{\eta}_v \}.
\]
with mutually distinct \( \alpha_i, \beta_j, \sigma_i, \eta_i \in \{1, \ldots, n\} \). Let \( \sigma = \{\sigma_1, \ldots, \sigma_u\}, \eta = \{\eta_1, \ldots, \eta_v\} \). Note that

\[
v = \begin{cases} u + 2, & J \in \mathcal{J}^{n-2}, \\ u + 1, & J \in \mathcal{J}^{n-1}. \end{cases}
\]

Let \( J \in \mathcal{J}^n \) be increasing. We introduce the following conditions for \( \varepsilon, \varepsilon' \in E_n \):

**Condition \( D_{J,\varsigma} \):**

- \( w_1(\varepsilon) + w_1(\varepsilon') = w_2(J) \),
- if \( w_1(\varepsilon) - w_1(\varepsilon') = \sum_{k=1}^s \epsilon_{\gamma_k} - \sum_{k=1}^t \epsilon_{\delta_k} \), then \( \sigma \subseteq \gamma \), and \( t \equiv r + \varsigma \pmod{2} \).

Here \( \varsigma \in \{0, 1\} \).

**Condition \( D_J \):**

- \( w_1(\varepsilon) + w_1(\varepsilon') = w_2(J) \),
- if \( w_1(\varepsilon) - w_1(\varepsilon') = \sum_{k=1}^s \epsilon_{\gamma_k} - \sum_{k=1}^t \epsilon_{\delta_k} \), then \( \sigma \subseteq \gamma \), and \( t \equiv r \pmod{2} \).

Here \( w_1, w_2 \) are the weight functions defined by (4.17), (4.18). We set \( \gamma = \{\gamma_1, \ldots, \gamma_s\} \), \( \delta = \{\delta_1, \ldots, \delta_t\} \). It is easy to see \( \sigma \cup \eta = \gamma \cup \delta \).

In the following Conjectures 5.10, 5.11, \( g = D_n^{(1)} \), and \( \sigma \) and \( \eta \) are disjoint subsets of \( \{1, \ldots, n\} \).

**Conjecture 5.10.** Let \( J \) be an increasing element of \( \mathcal{J}^{n-2} \), and let \( \varsigma = 0, 1 \). Then we have

\[
(-1)^{\frac{u(u-1)}{2}} Q_{J,z}^{(n-2)} x_{\sigma}^{1-2u} x_{\eta}^{1-u} \prod_{i,j \neq i,j} (-x_i^{-1}) \prod_{1 \leq k \leq u, 1 \leq l \leq u+2, \sigma_k \not\in \eta_l} \left( \frac{-x_{\sigma_k}}{x_{\eta_l}} \right) \]

\[
= \sum (-1)^{\frac{s(s+1)}{2}} R_{\varepsilon,q}^{(n-\varsigma)} R_{\varepsilon',q-1}^{(n-\varsigma)} x_{\gamma}^{-2u} x_{\delta}^{-3-t} \prod_{1 \leq k \leq u, 1 \leq l \leq t, \eta_k \not\in \delta_l} \left( \frac{-x_{\gamma_k}}{x_{\delta_l}} \right),
\]

where the sum is taken over all \( \varepsilon, \varepsilon' \in E_{n,\varsigma} \) satisfying Condition \( D_{J,\varsigma} \), where \( \gamma \) and \( \delta \) are determined there.

**Conjecture 5.11.** Let \( J \) be an increasing element of \( \mathcal{J}^{n-1} \). Then we have

\[
(-1)^{\frac{u(u-1)}{2}} Q_{J,z}^{(n-1)} x_{\sigma}^{1-2u-u} x_{\eta}^{1-u} \prod_{i,j \neq i,j} (-x_i^{-1}) \prod_{1 \leq k \leq u, 1 \leq l \leq u+1, \sigma_k \not\in \eta_l} \left( \frac{-x_{\sigma_k}}{x_{\eta_l}} \right) \]

\[
= \sum (-1)^{\frac{t(t-1)}{2}} R_{\varepsilon,z}^{(n-1)} R_{\varepsilon',z-1}^{(n)} x_{\gamma}^{-2u} x_{\delta}^{-1-t} \prod_{1 \leq k \leq u, 1 \leq l \leq t, \eta_k \not\in \delta_l} \left( \frac{-x_{\gamma_k}}{x_{\delta_l}} \right),
\]

where the sum is taken over all \( \varepsilon \in E_{n,1} \), \( \varepsilon' \in E_{n,0} \) satisfying Condition \( D_J \), where \( \gamma \) and \( \delta \) are determined there.

For \( n = 4 \), we have checked by Mathematica 5.0 the conjectures hold up to certain degree. For general \( n \), we prove in Appendix D the conjectures specialized to the ordinary characters \( \bar{Q}_{J}^{(a)} \) and \( \bar{R}_{\varepsilon}^{(n-1)}, \bar{R}_{\varepsilon}^{(n)} \).
6 Conclusion

In this section, we give a summary of this paper. We have done the following things:

1. To each weight of the fundamental representation $V^{(a)}$, we associated a formal series $Q_{J,z}^{(a)}$ or $R_{z}^{(a)}$. We expect that with some simple factors the formal series are $q$-characters of certain irreducible modules of $U_q(b)$.

2. Under suitable identifications, using relations for the connection coefficients implied by the $\psi$-system, we proposed the following conjecture relations for the series $Q_{J,z}^{(a)}$ and $R_{z}^{(a)}$:

   - $g = A_n^{(1)}$, Proposition 5.1. This is the Plücker type relations which is in fact the defining relations of the series $Q_{J,z}^{(a)}$;
   - $g = A_n^{(1)}$, Theorem 5.2. This is the Wronskian identity for the $q$-characters of $U_q(b)$ to be proved in Appendix B.
   - $g = B_n^{(1)}$, Conjecture 5.3.
   - $g = C_n^{(1)}$, Conjectures 5.5–5.8.
   - $g = D_n^{(1)}$, Conjectures 5.10, 5.11.

For the last three cases, we support our conjectures by checking the following. For the special cases $g = B_2^{(1)}$, $g = C_2^{(1)}$, we proved the conjectures by direct computations. For the cases $g = B_3^{(1)}$ and $g = D_4^{(1)}$, we checked the conjectures up to some degrees by Mathematica. When specialized to characters, the conjectured relations hold in all cases. This will be proved in Appendix C or D.

We hope we have presented reasonable grounds to suggest that the correspondence between the connection coefficients of certain differential equations and the $q$-characters of the Borel subalgebra $U_q(b)$ supplies an effective way to find polynomial relations. Certainly this is only the first step, and more serious checks are desirable along with attempts toward proving these identities.

A Vector representation of $L_{\mathfrak{g}}$

Following [14, Appendix 2], we give an explicit realization of the vector representation of $L_{\mathfrak{g}}$. The symbol $E_{i,j}$ stands for the matrix unit $(\delta_{i,a}\delta_{j,b})_{a,b=1,\ldots,N}$ where $N = \dim V^{(1)}$.

$L_{\mathfrak{g}} = A_n^{(1)}$:

\[ e_0 = E_{n+1,1}, \quad e_i = E_{i,i+1}, \quad 1 \leq i \leq n, \]
\[ f_0 = E_{1,n+1}, \quad f_i = E_{i+1,i}, \quad 1 \leq i \leq n, \]
\[ h_0 = -E_{1,1} + E_{n+1,n+1}, \quad h_i = E_{i,i} - E_{i-1,i+1}, \quad 1 \leq i \leq n, \]
\[ h_\rho = \text{diag} \left( -\frac{n}{2}, \ldots, \frac{n}{2} \right), \quad \ell = \text{diag} (\mu_1, \ldots, \mu_{n+1}), \quad \sum_{i=1}^{n+1} \mu_i = 0. \]

$L_{\mathfrak{g}} = A_{2n-1}^{(2)}$:

\[ e_0 = \frac{1}{2}(E_{1,2n-1} + E_{2,2n}), \quad e_i = E_{i+1,i} + E_{2n+1-i,2n-i}, \quad 1 \leq i \leq n - 1, \quad e_n = E_{n+1,n}, \]
\[ f_0 = 2(E_{2n-1,1} + E_{2n,2}), \quad f_i = E_{i,i+1} + E_{2n-i,2n+1-i}, \quad 1 \leq i \leq n - 1, \quad f_n = E_{n,n+1}, \]
\[ h_0 = E_{1,1} + E_{2,2} - E_{2n-1,2n-1} - E_{2n,2n}, \]
\[ h_i = -E_{i,i} + E_{i+1,i+1} - E_{2n-i,2n-i} + E_{2n+1-i,2n+1-i}, \quad 1 \leq i \leq n-1, \]
\[ h_{n} = -E_{n,n} + E_{n+1,n+1}, \]
\[ h_{\rho} = \text{diag}\left(-\frac{1}{2}, \ldots, -\frac{1}{2}, \frac{1}{2}, \ldots, n - \frac{1}{2}\right), \quad \ell = \text{diag}(\mu_1, \ldots, \mu_n, -\mu_n, \ldots, -\mu_1). \]

**Proof of Theorem 5.2**

In this section we give a proof of Theorem 5.2. Let \( g = A_n^{(1)} \), and set

\[ A_{i,z} = Y_{i,q^{-1}}^{(1)} Y_{i,q}^{-1} Y_{i-1,z} Y_{i+1,z}, \quad i = 1, \ldots, n. \]

One can rewrite \( Q_{i,z}^{(1)} \) as

\[ Q_{i,z}^{(1)} = Y_{n,q^{-1}z}^{-1} \sum_{k_{i+1}, \ldots, k_{n+1}} \prod_{j=1}^{n} \left( \frac{x_{j+1}}{x_i} \right)^{k_{j+1}} A_{j,q^{-2}z}^{-1} \sum_{j=1}^{n+1} k_{j+1} \right). \]

For \( i \in J \), we define

\[ \Delta_{i,z}^{(0)} = Y_{n,q^{-1}z} Q_{i,z}^{(1)}, \]
\[ \Delta_{i,z}^{(a)} = A_{n-i,q^{-1}z}^{-1} \left( \Delta_{i,z}^{(a-1)} - \frac{x_{n+2-a}}{x_i} \right), \quad a = 1, \ldots, n. \]

By the above definition and the formula of \( Q_{i,z}^{(1)} \), we have the following lemma.
Lemma B.1. For \( i \in \mathcal{J} \) and \( 0 \leq a \leq n \), we have

\[
\Delta_{i,z}^{(a)} = \begin{cases} 
\sum_{k_{i+1}, \ldots, k_{n+1} = i}^{n-a} \left( \frac{x_{j+1}}{x_i} \right)^{k_{j+1}} A_{n-2}^{-1} j q^{1-a}_{j-\nu} z \right) \quad & a < n + 1 - i, \\
1, \quad & a = n + 1 - i, \\
0, \quad & a > n + 1 - i.
\end{cases}
\]

Proposition B.2. For an element \( J = (j_1, j_2, \ldots, j_a) \in \mathcal{J}^a \), we have

\[
Q_{J,z}^{(a)} = \sum_{n+1-a,q^{n+1}z}^{n+1-a,q^{n+1}z} \left( \frac{a}{x_i} \right)^{k_{j+1}} A_{n-2}^{-1} j q^{1-a}_{j-\nu} z \right).
\]

Proof. From the definition we have

\[
Q_{J,z}^{(a)} = \prod_{k=1}^{a} n \cdot q^{n-a+2k} z
\]

For \( k = 2, \ldots, a \), we subtract the \((k-1)\)-th row multiplied by \( z_{n+1} \) from the \( k\)-th row and then extract the factor \( A_{n,q^{n-a-1+2k}z}^{-1} \). Thus we get

\[
Q_{J,z}^{(a)} = \prod_{k=2}^{a} n \cdot q^{n-a+2k} z \prod_{k=2}^{a} A_{n,q^{n-a-1+2k}z}^{-1}
\]

Repeating the above steps in the order \( b = 2, \ldots, a - 1 \), by replacing \( x_{n+1} \) and \( A_{n,q^{n-a-1+2k}z}^{-1} \) with \( x_{n+2-b} \) and \( A_{n+1-b,q^{n-a-b+2k}z}^{-1} \), we get

\[
Q_{J,z}^{(a)} = \prod_{b=1}^{a} n \cdot q^{n-a+2k} z \prod_{b=1}^{a-1} b+1 \prod_{b=1}^{a} A_{n+1-b,q^{n-a-b+2k}z}^{-1}
\]

The proof is over by noting that

\[
\prod_{k=1}^{a} n \cdot q^{n-a+2k} z \prod_{b=1}^{a} b+1 \prod_{b=1}^{a} A_{n+1-b,q^{n-a-b+2k}z}^{-1} = \gamma_{n+1-a,q^{n+1}z}^{-1}.
\]

■
Now we give a proof of Theorem 5.2.

**Proof of Theorem 5.2.** By Proposition B.2, for \( J = (1, 2, \ldots, n + 1) \), we have

\[
Q_{J,z}^{(n+1)} = \begin{vmatrix}
\Delta^{(0)}_{1,q^{-n}z} x_1^0 & \Delta^{(0)}_{2,q^{-n}z} x_2^0 & \cdots & \Delta^{(0)}_{n+1,q^{-n}z} x_{n+1}^0 \\
\Delta^{(1)}_{1,q^{-n}z} x_1 & \Delta^{(1)}_{2,q^{-n}z} x_2 & \cdots & \Delta^{(1)}_{n+1,q^{-n}z} x_{n+1} \\
\vdots & \vdots & \ddots & \vdots \\
\Delta^{(n)}_{1,q^n z} x_1^n & \Delta^{(n)}_{2,q^n z} x_2^n & \cdots & \Delta^{(n)}_{n+1,q^n z} x_{n+1}^n \\
\end{vmatrix}.
\]

By Lemma B.1, we get

\[
Q_{J,z}^{(n+1)} = (-1)^{(n+1)n/2} \prod_{b=1}^{n} x_b^{n+1-b},
\]

which implies the result. 

\[\blacksquare\]

## C Proofs of identities for characters in the case \( B_n^{(1)} \)

In this section, we set \( g = B_n^{(1)} \). First we introduce some notation. For \( i, j \in J \) and an element \( J \in J^n \), we define

\[
h_J = \prod_{k, k \in J} \left(x_k^{1/2} + x_k^{-1/2}\right), \quad \bar{x}_{i,j} = \prod_{k \in J, k < i} \left(-x_k^{1/2} b_{k,i}\right) \prod_{k \in J^*, \iota < k} \left(x_k^{1/2} b_{k,i}\right),
\]

\[
g_{i,j} = f_{i,j} x_j^{1/2} b_{j,i} = -g_{j,i}, \quad g_{i,J} = \prod_{k \in J, \iota \neq i} g_{i,k},
\]

where \( f_{i,j} \) is defined in Subsection 4.2.

By (4.15) and the above notation, one has

\[
\bar{Q}_J^{(a)} = \prod_{i \in J, j \in J^*} f_{i,j} \prod_{i < j} ((x_j - x_i) f_{i,j}) = X_J^{-1} h_J \prod_{i \in J, j \in J^*} f_{i,j},
\]

where \( X_J \) is given in (5.5). A direct computation gives

\[
\prod_{i \in J, j \in J^*} f_{i,j} = \prod_{i \in J, j \in J^*} f_{i,j}, \quad \prod_{j \in J, \iota < j} \bar{x}_{i,j} = \prod_{j \in J, \iota < j} \bar{x}_{i,j},
\]

\[
\prod_{i \in J - \{j\}, j \in J^*} f_{i,j} = \prod_{i \in J - \{j\}, j \in J^*} f_{i,j}, \quad \prod_{j \in J, \iota < j} g_{j,k} \quad \bar{x}_{j,k} = \prod_{j \in J^*} g_{j,k} \bar{x}_{j,k}.
\]

Furthermore, for two elements \( J_1 \) and \( J_2 \) of \( J^n \) and \( j_k \in J_2 \cap J_1^* \), we have

\[
\frac{x_{j_k,l} \bar{x}_{j_k,l}^2}{x_{j_k,l} \bar{x}_{j_k,l}} = -x_{j_k}^2 x_{j_1} x_{j_2}^{-1}.
\]

Using the above formulas, the specialization of Conjecture 5.3 to the characters \( \bar{Q}_J^{(a)} \) reduces to the following proposition.
Proposition C.1. For two increasing elements $J_1 = (i_1, \ldots, i_n)$, $J_2 = (j_1, \ldots, j_n) \in J^n$, we have

$$x_{J_1}^{-\frac{1}{2}} h_{J_1} h_{J_2} - x_{J_2}^{-\frac{1}{2}} h_{J_2} h_{J_1} = x_{J_1}^{-1} x_{J_2}^{-1} \sum_{j \in J_2 \cap J_1^c} x_{J_1}^{-\frac{3}{2}} x_{J_2}^{-\frac{3}{2}} g_{J_1,j}^* g_{J_2,j} h_{J_2-j} h_{J_1-j}.$$  \(\text{(C.1)}\)

Proof. Generally let

$$T_1 = \left\{ j \in J_1 \cap J_2^* \mid \tilde{j} \notin J_1, \tilde{j} \notin J_2^* \right\},$$

$$T_2 = \left\{ j \in J_1 \cap J_2^* \mid \tilde{j} \notin J_1, \tilde{j} \notin J_2^* \right\}, \quad T_3 = \left\{ j \in J_1 \cap J_2^* \mid \tilde{j} \notin J_1, \tilde{j} \notin J_2^* \right\},$$

$$S_1 = \left\{ \tilde{j} \notin J_1, \tilde{j} \notin J_2^* \right\}, \quad S_2 = \left\{ \tilde{j} \notin J_2, \tilde{j} \notin J_1^* \right\}, \quad S_3 = \left\{ \tilde{j} \notin J_2, \tilde{j} \notin J_1^* \right\},$$

then we have $J_1 \cap J_2^* = \bigcup_{1 \leq i \leq 3} T_i \cup T_3$, $J_2 \cap J_1^* = \bigcup_{1 \leq i \leq 3} S_i \cup S_3$, and $\#(J_1 \cap J_2^*) = \#(J_2 \cap J_1^*)$. Here for a subset $S \subseteq J$, we set $\mathcal{S} = \{ i \mid i \in S \}$.

Let

$$F_{J_1,J_2}(x_j) = x_j^3 x_{J_1}^{-1} h_{J_2-j} h_{J_1-j}^* g_{J_1,j} g_{J_2,j}.$$

We also introduce a function

$$f(z) = \frac{1}{(1+z)} \prod_{k \in T_1} (z-x_k) \prod_{k \in S_1} (z-x_k) \prod_{k \in T_3} \frac{x_k-1}{z-x_k} \prod_{k \in S_3} \frac{(z-x_k)(z x_k - 1)}{(z-x_k)(z x_k - 1)},$$

which has poles only at $z = -1, \infty$ and $x_j \ (j \in J_2 \cap J_1^* = \bigcup_{1 \leq i \leq 3} S_i \cup S_3)$.

A direct computation gives, for $j \in S_3$

$$\text{Res}(f, x_j^{-1}) = \frac{1}{x_j(1 + x_j)(1 - x_j)} \prod_{k \in T_1} (1 - x_k x_j) \prod_{k \in S_1} (1 - x_k x_j) \prod_{k \in T_3} \frac{x_k - x_j}{1 - x_k x_j} \prod_{k \in S_3, \ k \neq j} \frac{(x_j - x_k)(x_j x_k - 1)}{(x_j - x_k)(x_j x_k - 1)} = c F_{J_1,J_2}(x_j^{-1}),$$

where $c = \frac{x_{T_3} x_{S_2}}{x_{S_3}}$. Similarly we obtain

$$\text{Res}(f, x_j) = c F_{J_1,J_2}(x_j), \quad j \in S_i, \quad i = 1, 2, 3.$$

On the other hand, by setting $h_i = x_{\frac{i}{2}} + x_{\frac{i}{2}}^{-1}$, it is easy to see that

$$\text{Res}(f, -1) = \frac{x_{T_3}}{x_{S_3}} \frac{x_{T_1}}{x_{S_1}} \prod_{k \in T_1} h_k \prod_{k \in T_3} h_k^2 \prod_{k \in S_1} h_k \prod_{k \in S_3} h_k^2, \quad \text{Res}(f, \infty) = -\frac{x_{T_3} x_{S_2}}{x_{S_3}} = -c.$$
Since the right-hand side of (C.1) equals \( h_{j_1} h_{j_2} x_{j_1}^{-\frac{1}{2}} \sum_{j \in J_2 \cap J_1^c} F_{j_1, j_2}(x_j) \), it simplifies to

\[
h_{j_1} h_{j_2} x_{j_1}^{-\frac{1}{2}} (-c^{-1}) \left( -c + \frac{x_T}{x_{S_3}} \left( \frac{x_{T_1}}{x_{S_1}} \right)^{\frac{1}{2}} \prod_{k \in T_1} h_k \prod_{k \in T_3} h_k^2 \right).
\]

Now it is sufficient to prove

\[
\left( \frac{x_{T_1}}{x_{S_1}} \right)^{\frac{1}{2}} \prod_{k \in T_1} h_k \prod_{k \in T_3} h_k^2 = x_{S_2} \frac{1}{x_{j_1}} h_{j_1} h_{j_2}
\]

which is immediate. We finish the proof.

\[
\text{D Proofs for identities of characters of case } C_n^{(1)} \text{ and } D_n^{(1)}
\]

In this section we consider \( g = C_n^{(1)}, D_n^{(1)} \). By (4.15) and Conjecture 4.5, Conjectures 5.5–5.8 reduce to the following identities

\[
(-1)^{\frac{u(u+1)}{2}} x_{\sigma} \prod_{1 \leq k \leq u, 1 \leq s \leq u+1} \frac{1}{1 - x_{\sigma_k} x_{\eta_l}} \prod_{1 \leq k \leq u+1} (1 - x_{\eta_k}^2)
\]

\[
= \sum (-1)^{\frac{s(s+1)}{2}} x_{\gamma} \prod_{1 \leq k \leq s} \frac{1}{1 - x_{\gamma_k} x_{\eta_l}} \prod_{1 \leq k \leq s+1} (1 - x_{\eta_k}) \prod_{1 \leq k \leq s} (1 - x_{\eta_k}^2).
\]

\[
(-1)^{\frac{u(u+1)}{2}} x_{\sigma} \prod_{1 \leq k \leq u, 1 \leq s \leq u+2} \frac{1}{1 - x_{\sigma_k} x_{\eta_l}} \prod_{1 \leq k \leq u+2} (1 - x_{\eta_k}^2)
\]

\[
= \sum (-1)^{\frac{s(s+1)}{2}} x_{\gamma} \prod_{1 \leq k \leq s} \frac{1}{1 - x_{\gamma_k} x_{\eta_l}} \prod_{1 \leq k \leq s+2} (1 - x_{\eta_k}) \prod_{1 \leq k \leq s} (1 - x_{\eta_k}^2).
\]

\[
(-1)^{\frac{u(u+1)}{2}} x_{\sigma} \prod_{1 \leq k \leq u, 1 \leq s \leq u+1} \frac{1}{1 - x_{\sigma_k} x_{\eta_l}} \prod_{1 \leq k \leq u+1} (1 - x_{\eta_k}^2)
\]

\[
= \sum (-1)^{\frac{s(s+1)}{2}} x_{\gamma} \prod_{1 \leq k \leq s} \frac{1}{1 - x_{\gamma_k} x_{\eta_l}} \prod_{1 \leq k \leq s+1} (1 - x_{\eta_k}) \prod_{1 \leq k \leq s} (1 - x_{\eta_k}^2).
\]

In these formulas, \( \sigma, \eta \) are disjoint subsets of \( \{1, \ldots, n\} \), such that \( (\sigma, \eta) = (u, u+1), (u, u+2), (u, u), (u, u+1) \) respectively. The sum is taken over partitions of \( \sigma \cup \eta \) into subsets \( \gamma, \delta \) satisfying \( \sigma \subseteq \gamma \), where we set \( s = \frac{1}{2} \gamma, t = \frac{1}{2} \delta \).

Similarly, in the case of \( g = D_n^{(1)} \), by (4.15) and Conjecture 4.6 the relevant identities read as follows.

\[
(-1)^{\frac{u(u+1)}{2}} x_{\sigma} \prod_{1 \leq k \leq u, 1 \leq s \leq u+2} \frac{1}{1 - x_{\sigma_k} x_{\eta_l}} \prod_{1 \leq k \leq u+2} (1 - x_{\eta_k}^2)
\]

\[
= \sum (-1)^{\frac{s(s+1)}{2}} x_{\gamma} \prod_{1 \leq k \leq s} \frac{1}{1 - x_{\gamma_k} x_{\eta_l}} \prod_{1 \leq k \leq s+2} (1 - x_{\eta_k}) \prod_{1 \leq k \leq s} (1 - x_{\eta_k}^2).
\]
where \( \tau, \tau' \in \mathbb{Z}, \), \( u, u + 2 \), \( u, u + 1 \) respectively, \( s = \tau \gamma, t = \tau' \delta, \) and the sum is taken over the partitions satisfying \( \sigma \subseteq \gamma \) keeping fixed the parity of \( s \).

These formulas are related by specialization of variables: (D.1) is obtained from (D.3) by setting \( x_{\sigma_u} = 0, \) (D.2) is obtained from (D.4) by setting \( x_{\sigma_u} = 0, \) and (D.4) is obtained from (D.3) by setting \( x_{\sigma_u} = 1. \) Likewise (D.5) is obtained from (D.6) by setting \( x_{\sigma_u} = \infty. \) Hence it is sufficient to deal only with (D.3) and (D.6).

Define polynomials in \( x = (x_1, \ldots, x_m) \) and \( y = (y_1, \ldots, y_n), \)

\[
F^{(m,n)}(x,y) = \prod_{1 \leq i < j \leq m} (1 - x_ix_j) \prod_{1 \leq i < j \leq n} (1 - y_iy_j)(y_j - y_i),
\]

\[
G^{(m,n)}_{J_1,J_2}(x,y) = \prod_{1 \leq i \leq m, j \in J_1} (1 - x_iy_j) \prod_{1 \leq j \leq n, i \in J_2} (x_i - y_j)
\]

\[
\times \prod_{i,j \in J_1 \backslash J_2} (y_j - y_i) \prod_{i \in J_1, j \in J_2} (1 - y_iy_j) \prod_{i,j \in J_2 \backslash J_1} (y_i - y_j).
\]

Set also \( \delta(J_1, J_2) = \tau \{ (i, j) \in J_1 \times J_2 \mid i < j \}. \) Then (D.3) and (D.6) can be rewritten respectively as follows:

\[
\prod_{j=1}^n (1 - y_j^2) \cdot F^{(n,n)}(x,y) = \sum_{J_1,J_2} (-1)^{\delta(J_1,J_2)} \prod_{j \in J_2} y_j \cdot G^{(n,n)}_{J_1,J_2}(x,y); \quad (D.7)
\]

\[
(1 + \epsilon) \prod_{j=1}^{n-1} (1 - x_j^2) \cdot F^{(n-1,n)}(x,y) = \sum_{J_1,J_2} \epsilon^{\tau J_2} (-1)^{\delta(J_1,J_2)} G^{(n-1,n)}_{J_1,J_2}(x,y). \quad (D.8)
\]

Here the sum in the right-hand sides are taken over all partitions \( J_1 \sqcup J_2 \) of \( \{1, \ldots, n\}, \) and \( \epsilon = \pm 1. \) For the second, we have added/subtracted the original sums which have a parity restriction.

In order to show (D.7), we prove it in a slightly more general form.

**Proposition D.1.** Suppose \( n \geq 2, \) \( 0 \leq m \leq n \) and \( m \equiv n \mod 2. \) Then

\[
\delta_{m,n} \prod_{j=1}^n (1 - y_j^2) \cdot F^{(m,n)}(x,y) = \sum_{J_1,J_2} (-1)^{\delta(J_1,J_2)} \prod_{j \in J_2} y_j \cdot G^{(m,n)}_{J_1,J_2}(x,y). \quad (D.9)
\]

**Proof.** Denote by \( R^{(m,n)} \) the right-hand side of (D.9), and set

\[
L^{(m,n)} = \prod_{j=1}^n (1 - y_j^2) \cdot F^{(m,n)}(x,y), \quad h^{(m)}(x,z) = \prod_{i=1}^m (1 - x_i z).
\]

Let us rewrite \( R^{(m,n)} \) as follows:

\[
\prod_{j=1}^n y_j^{-(m+n)/2} \cdot R^{(m,n)} = \sum_{J_1,J_2} (-1)^{\delta J_2} \prod_{i \in J_1} \left( y_i^{-(m+n)/2} h^{(m)}(x,y_i) \right) \prod_{j \in J_2} \left( y_j^{(m+n)/2} h^{(m)}(x,y_j^{-1}) \right).
\]

This completes the proof.
\[
\times \prod_{i<j}^{i,j \in J_1, i \neq j} (y_j - y_i) \prod_{i,j \in J_2, i > j} (y_j - y_i) \prod_{i,j \in J_2, i < j} (y_j - y_i)
\]
\[
= \sum_{\epsilon_1, \ldots, \epsilon_n = \pm 1} \prod_{i,j} (\epsilon_j y_j - \epsilon_i y_i)^{\frac{m+n}{2}} h(m)(x, y_i)
\prod_{1 \leq i < j \leq n} (y_j - y_i)
\]
\[
= \det \left( y_i^{-(m+n)/2} - y_{i-1}^{-(m+n)/2} h(m)(x, y_i - y_i) \right)_{1 \leq i, j \leq n}.
\]

It is easy to see that the last expression is 0 if \(m = 0\) or \(m = 1\). Assume \(m \geq 2\). We prove (D.9) by induction on \(m\). From the skew-symmetry of the determinant, \(R^{(m,n)}\) is divisible by \(\prod_{j=1}^{n} (1 - y_j^2) \prod_{1 \leq i < j \leq n} (1 - y_i y_j)(y_j - y_i)\). Let us show that it is divisible also by \(\prod_{1 \leq i < j \leq m} (1 - x_i x_j)\).

If we set \(x_m = x_{m-1}\), then \(h(m)(x, z) = z(z + 1 - x_{m-1} - x_{m-1}) h(m-2)(x', z)\), where \(x' = (x_1, \ldots, x_{m-2})\). Hence \(R^{(m,n)}\) becomes proportional to
\[
\det \left( y_i^{-(m-2+n)/2} - y_{i-1}^{-(m-2+n)/2} h(m-2)(x', y_i - y_i) \right)_{1 \leq i, j \leq n},
\]
which is 0 by the induction hypothesis. This shows that \(R^{(m,n)}\) is divisible by \(L^{(m,n)}\).

The total degree of \(L^{(m,n)}\) is \(l^{(m,n)} = mn - 1 + n(3n - 1)/2\), while that of \(R^{(m,n)}\) is at most \(r^{(m,n)} = n^2/2 + mn + (m-1)^2/4 - m/2\). Notice that \(l^{(m,n)} > r^{(m,n)}\) for \(n > m\) and \(l^{(n,n)} = r^{(n,n)}\). From this we conclude that if \(n > m\) then \(R^{(m,n)} = 0\), and if \(n = m\) then \(R^{(n,n)}\) is a constant multiple of \(L^{(n,n)}\). The constant is shown to be 1 by setting \(x = 0\) and using the Weyl denominator formula of type \(C_n\),
\[
\prod_{j=1}^{n} y_j^2 \det \left( y_i^{-(n+j-1)} - y_{i-1}^{-(n+j-1)} \right)_{1 \leq i, j \leq n} = \prod_{j=1}^{n} (1 - y_j^2) \prod_{1 \leq i < j \leq n} (1 - y_i y_j)(y_j - y_i).
\]

Proposition D.2. Suppose \(n \geq 3\) and \(0 \leq m \leq n - 1\). Then for \(\epsilon = \pm 1\) we have
\[
(1 + \epsilon) \delta_{m,n-1} \prod_{i=1}^{m} (1 - x_i^2) \cdot F^{(m,n)}(x, y) = \sum_{J_1, J_2} \epsilon^{\delta_{J_2}}(-1)^{\delta_{J_1, J_2}} G_{J_1, J_2}^{(m,n)}(x, y).
\]

Proof. We denote the right hand-side of (D.10) by \(R^{(m,n)}_{\epsilon}\), and set
\[
K^{(m,n)} = \prod_{i=1}^{m} (1 - x_i^2) \cdot F^{(m,n)}(x, y).
\]

By a calculation similar to the one in the proof of Proposition D.1 we find
\[
\prod_{j=1}^{n} y_j^{-(m+n+1)/2} R^{(m,n)}_{\epsilon}
\]
\[
= \det \left( y_i^{-(m+n+1)/2} + \epsilon' y_i^{-(m+n+1)/2} h(m)(x, y_i) + \epsilon' y_i^{-(m+n+1)/2} h(m)(x, y_i^{-1}) \right)_{1 \leq i, j \leq n},
\]
where \(\epsilon' = (-1)^{m+n-1} \epsilon\). From this, it is easy to see that \(R^{(0,n)}_{\epsilon} = 0\). As in Proposition D.1, we prove the assertion by induction on \(m\). As before, one can verify that the left-hand side is divisible by \(\prod_{1 \leq i < j \leq m} (1 - x_i x_j) \prod_{1 \leq i < j \leq n} (1 - y_i y_j)(y_j - y_i)\). Let us show that it is divisible also by \(\prod_{i=1}^{m} (1 - x_i^2)\). If we set \(x_m = \pm 1\), then we have
\[
h(m)(x, z) \bigg|_{x_m = \pm 1} = (z^{-1/2} + z^{1/2}) \cdot z^{1/2} h(m-1)(x', z), \quad x' = (x_1, \ldots, x_{m-1}).
\]
Hence the determinant becomes proportional to
\[ \det \left( y_i^{-\lfloor (m+n)/2+j \rfloor h^{(m-1)}(x, y_i) + \epsilon y_i^{(m+n)/2-j h^{(m-1)}(x, y_i)} \right)_{1 \leq i, j \leq n}, \]
which is 0 by the induction hypothesis.

The total degree of \( K^{(m,n)} \) is \( k^{(m,n)} = m(m+1) + 3n(n-1)/2 \), while that of \( R^{(m,n)}_\epsilon \) is at most
\[ r^{(m,n)}_\epsilon = mn + n(n-1)/2 + (m+n)^2/4 - \delta/4, \]
where \( \delta = 0, 1 \) is determined by \( \delta \equiv m-n \pmod{2} \). Then \( k^{(m,n)} > r^{(m,n)}_\epsilon \) if \( m \leq n - 2 \), and \( k^{(n-1,n)} = r^{(n-1,n)}_\epsilon \). Hence we find that \( R^{(m,n)}_\epsilon = 0 \) if \( m \leq n - 2 \), and \( R^{(n-1,n)}_\epsilon \) is a constant multiple of \( K^{(n-1,n)} \). The constant can be found by setting \( x = 0 \) and using the Weyl denominator formula of type \( D_n \),
\[ \prod_{j=1}^n y_j^n \cdot \det \left( y_i^{-n+j} + \epsilon y_i^{n-j} \right)_{1 \leq i, j \leq n} = (1 + \epsilon) \prod_{1 \leq i < j \leq n} (1 - y_i y_j)(y_j - y_i). \]

**Acknowledgements**

The author would like to express her deep gratitude to Professor Michio Jimbo for his valuable advices, great helps and kindness, and to Professor Hidetaka Sakai for his constant encouragement during the course of the present work. She wishes to thank also Professor Junji Suzuki for helpful discussions. Finally the author would like to thank the anonymous referees for comments and suggestions that improved the paper greatly.

**References**

[1] Baxter R.J., Partition function of the eight-vertex lattice model, *Ann. Physics* 70 (1972), 193–228.
[2] Bazhanov V.V., Frassek R., Lukowski T., Meneghelli C., Staudacher M., Baxter Q-operators and representations of Yangians, *Nuclear Phys. B* 850 (2011), 148–174, arXiv:1010.3699.
[3] Bazhanov V.V., Hibberd A.N., Khoroshkin S.M., Integrable structure of \( \mathcal{W}_3 \) conformal field theory, quantum Boussinesq theory and boundary affine Toda theory, *Nuclear Phys. B* 622 (2002), 475–547, hep-th/0105177.
[4] Bazhanov V.V., Lukyanov S.L., Zamolodchikov A.B., Integrable structure of conformal field theory. II. Q-operator and DDV equation, *Comm. Math. Phys.* 190 (1997), 247–278, hep-th/9604044.
[5] Bazhanov V.V., Tsutsumi Z., Baxter’s Q-operators for supersymmetric spin chains, *Nuclear Phys. B* 805 (2008), 451–516, arXiv:0805.4274.
[6] Beck J., Braid group action and quantum affine algebras, *Comm. Math. Phys.* 165 (1994), 555–568, hep-th/9404165.
[7] Bowman J., Irreducible modules for the quantum affine algebra \( U_q(g) \) and its Borel subalgebra \( U_q(g)^{\geq 0} \), *J. Algebra* 316 (2007), 231–253, math.QA/0606627.
[8] Chari V., Greenstein J., Filtrations and completions of certain positive level modules of affine algebras, *Adv. Math.* 194 (2005), 296–331, math.QA/0309016.
[9] Chari V., Pressley A., Quantum affine algebras, *Comm. Math. Phys.* 142 (1991), 261–283.
[10] Chari V., Pressley A., Quantum affine algebras and their representations, in Representations of Groups (Banff, AB, 1994), *CMS Conf. Proc.*, Vol. 16, Amer. Math. Soc., Providence, RI, 1995, 59–78, hep-th/9411145.
[11] Dorey P., Dunning C., Masoero D., Suzuki J., Tateo R., Pseudo-differential equations, and the Bethe ansatz for the classical Lie algebras, *Nuclear Phys. B* 772 (2007), 249–289, hep-th/0612298.
[12] Dorey P., Dunning C., Tateo R., The ODE/IM correspondence, *J. Phys. A: Math. Theor.* 40 (2007), R205–R283, hep-th/0703066.
[13] Drinfel’d V.G., A new realization of Yangians and of quantum affine algebras, *Sov. Math. Dokl.* 36 (1987), 212–216.
[14] Drinfel’d V.G., Sokolov V.V., Lie algebras and equations of Korteweg–de Vries type, *J. Math. Sci.* 30 (1985), 1975–2036.
[15] Feigin B., Frenkel E., Quantization of soliton systems and Langlands duality, in Exploring New Structures and Natural Constructions in Mathematical Physics, Adv. Stud. Pure Math., Vol. 61, Math. Soc. Japan, Tokyo, 2011, 185–274, arXiv:0705.2486.

[16] Frenkel E., Mukhin E., Combinatorics of $q$-characters of finite-dimensional representations of quantum affine algebras, Comm. Math. Phys. 216 (2001), 23–57, math.QA/9911112.

[17] Frenkel E., Reshetikhin N., The $q$-characters of representations of quantum affine algebras and deformations of $W$-algebras, in Recent Developments in Quantum Affine Algebras and Related Topics (Raleigh, NC, 1998), Contemp. Math., Vol. 248, Amer. Math. Soc., Providence, RI, 1999, 163–205, math.QA/9810055.

[18] Fulton W., Young tableaux. With applications to representation theory and geometry, London Mathematical Society Student Texts, Vol. 35, Cambridge University Press, Cambridge, 1997.

[19] Hernandez D., The Kirillov–Reshetikhin conjecture and solutions of $T$-systems, J. Reine Angew. Math. 596 (2006), 63–87, math.QA/0501202.

[20] Hernandez D., Jimbo M., Asymptotic representations and Drinfeld rational fractions, arXiv:1104.1891.

[21] Inoue R., Iyama O., Kuniba A., Nakanishi T., Suzuki J., Periodicities of $T$-systems and $Y$-systems, Nagoya Math. J. 197 (2010), 59–174, arXiv:0812.0667.

[22] Jantzen J.C., Lectures on quantum groups, Graduate Studies in Mathematics, Vol. 6, Amer. Math. Soc., Providence, RI, 1996.

[23] Kojima T., Baxter’s $Q$-operator for the $W$-algebra $W_N$, J. Phys. A: Math. Theor. 41 (2008), 355206, 16 pages, arXiv:0803.3505.

[24] Kuniba A., Nakanishi T., Suzuki J., Functional relations in solvable lattice models. I. Functional relations and representation theory, Internat. J. Modern Phys. A 9 (1994), 5215–5266, hep-th/9309137.

[25] Kuniba A., Nakanishi T., Suzuki J., $T$-systems and $Y$-systems in integrable systems, J. Phys. A: Math. Theor. 44 (2011), 103001, 146 pages, arXiv:1010.1344.

[26] Kuniba A., Suzuki J., Analytic Bethe ansatz for fundamental representations of Yangians, Comm. Math. Phys. 173 (1995), 225–264, hep-th/9406180.

[27] Nakai W., Nakanishi T., Paths, tableaux and $q$-characters of quantum affine algebras: the $C_n$ case, J. Phys. A: Math. Gen. 39 (2006), 2083–2115, math.QA/0502041.

[28] Nakajima H., $t$-analog of $q$-characters of Kirillov–Reshetikhin modules of quantum affine algebras, Represent. Theory 7 (2003), 259–274, math.QA/0204185.

[29] Tsuboi Z., Solutions of the $T$-system and Baxter equations for supersymmetric spin chains, Nuclear Phys. B 826 (2010), 399–455, arXiv:0906.2039.