Essential spectra of difference operators on \( \mathbb{Z}^n \)-periodic graphs

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Received 18 April 2007, in final form 10 July 2007
Published 1 August 2007
Online at stacks.iop.org/JPhysA/40/10109

Abstract
Let \((X, \rho)\) be a discrete metric space. We suppose that the group \(\mathbb{Z}^n\) acts freely on \(X\) and that the number of orbits of \(X\) with respect to this action is finite. Then we call \(X\) a \(\mathbb{Z}^n\)-periodic discrete metric space. We examine the Fredholm property and essential spectra of band-dominated operators on \(l^p(X)\) when \(1 < p < \infty\). Our approach is based on the theory of band-dominated operators on \(\mathbb{Z}^n\) and their limit operators. In the case where \(X\) is the set of vertices of a combinatorial graph, the graph structure defines a Schrödinger operator on \(l^p(X)\) in a natural way. We illustrate our approach by determining the essential spectra of Schrödinger operators with slowly oscillating potential both on zig-zag and on hexagonal graphs, the latter being related to nano-structures.

PACS numbers: 02.30.Tb, 03.65.Db
Mathematics Subject Classification: 81Q10, 46N50, 47B36

1. Introduction
In the past years, spectral properties of Schrödinger operators on quantum graphs have attracted a lot of attention due to their interesting mathematical properties and due to existing and expected applications in nano-structures as well (see, for instance, [4, 10, 40]). Quantum graph models also occur in chemistry and physics (see [28, 39] and [4, 16] and the references therein). The spectral properties of Schrödinger operators on quantum graphs were studied by P Kuchment and collaborators in a series of papers [15–20]. Direct and inverse spectral problems for Schrödinger operators on graphs connected with zig-zag carbon nano-tubes were considered in [13, 14]. Note that the structure of the spectra of periodic difference operators is essentially different from the continuous case. For instance, a periodic difference operator can possess eigenfunctions with compact support ([18, 20]), whereas this cannot happen
for a periodic differential operator of a second kind, which has an absolutely continuous spectrum.

It has been pointed out in [5, 27] and [17, 18] that the determination of the spectrum of a magnetic Laplacian on an equilateral quantum graph (i.e., a graph consisting of identical segments with the same potentials on them) can be reduced to the study of the spectrum of a discrete magnetic Laplacian. This observation makes difference operators on combinatorial graphs to an essential tool in the theory of differential operators on quantum graphs.

The main theme of this paper is the essential spectrum of difference operators (with the Schrödinger operators as a prominent example) acting on the spaces \( l^p(X) \) where \( X \) is the set of the vertices of a combinatorial graph \( \Gamma \). We exclusively consider discrete graphs \( \Gamma \) on which the group \( \mathbb{Z}^n \) acts freely and which have a finite fundamental domain with respect to this action.

For every \( 1 < p < \infty \), we introduce a Banach algebra \( A_p(X) \) of so-called band-dominated difference operators on \( l^p(X) \). Following [33, 34] and [35], we associate with every operator \( A \in A_p(X) \) a family \( \text{op}_p(A) \) of limit operators and show that an operator \( A \in A_p(X) \) is Fredholm on \( l^p(X) \) if and only if all operators in \( \text{op}_p(A) \) are invertible and if the norms of their inverses are uniformly bounded. In general, the limit operators of an operator \( A \) are simpler objects than the operator \( A \) itself. Thus, the limit operators method often provides an effective tool to study the Fredholm property of operators in \( A_p(X) \).

For operators in the so-called Wiener algebra \( \mathcal{W}(X) \) (which is a non-closed subalgebra of every algebra \( A_p(X) \)), the uniform boundedness of norms of inverse operators to limit operators follows already from their invertibility. This basic fact implies the useful identity

\[
\text{sp}_{\text{ess}} A = \bigcup_{A_h \in \text{op}_p A} \text{sp} A_h,
\]

where the set of the limit operators of \( A \), the spectra \( \text{sp} A_h \) of the limit operators of \( A \) and, hence, also the essential spectrum \( \text{sp}_{\text{ess}} A \) of \( A \), are independent of \( p \).

In the case of \( X = \mathbb{Z}^n \), formula (1) was obtained in [33, 35]. In [31], we applied this formula to electro-magnetic Schrödinger operators on the lattice \( \mathbb{Z}^n \). In particular, we determined the essential spectrum of the Hamiltonian of the 3-particle problem on \( \mathbb{Z}^n \).

In [29], one of the authors obtained an identity similar to (1) for perturbed pseudodifferential operators on \( \mathbb{R}^n \). He applied this result to study the location of the essential spectra of electro-magnetic Schrödinger, square-root Klein–Gordon, and Dirac operators under general assumptions with respect to the behavior of magnetic and electric potentials at infinity. On the basis of this method, he also gave a simple and transparent proof of the well-known Hunziker, van Winter, Zhislin theorem (HWZ-theorem) on the location of essential spectra of multi-particle Hamiltonians.

It should be noted that formulae similar to (1) have been obtained independently (but later) in [22] by means of admissible geometric methods. We also mention the papers [9, 8, 24, 3] and the references therein where \( C^* \)-algebra techniques were applied to study essential spectra of Schrödinger operators.

The present paper is organized as follows. In section 2 we collect some auxiliary material from [35] on matrix band-dominated operators on the lattice \( \mathbb{Z}^n \). In section 3 we introduce the Banach algebra \( A_p(X) \) of band-dominated operators acting on \( l^p(X) \) where \( X \) is a periodic discrete metric space on which the group \( \mathbb{Z}^n \) acts freely. We construct an isomorphism between the Banach algebra \( A_p(X) \) and the Banach algebra \( A_p(\mathbb{Z}^n, \mathbb{C}^N) \) of all (block) band-dominated operators on \( l^p(\mathbb{Z}^n, \mathbb{C}^N) \) where \( N \) is the number of points in the fundamental domain of \( X \) with respect to the action of \( \mathbb{Z}^n \). Applying this isomorphism and the results of section 2, we
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Derive necessary and sufficient conditions for an operator $A \in A_p(X)$ to be Fredholm. We also introduce a Wiener algebra $W(X)$ and derive formula (1) for operators in $W(X)$.

For later use, we recall some facts on periodic band-dominated operators in section 4. An operator $A \in A_p(X)$ is periodic if it commutes with each operator $L_h$ of left shift by $h \in \mathbb{Z}^n$. Periodic operators are distinguished by the fact that $\text{sp}_{\text{ess}} A = \text{sp} A$. With every periodic operator $A \in W(X)$, there is associated a continuous matrix function $\sigma_A$ called the symbol of $A$, which owns the property that the eigenvalues of $\sigma_A$ give the spectrum of $A$. In the terminology of [16, 17], $\sigma_A$ is the Floquet transform of $A$. We prefer to follow the theory of discrete convolutions and use the discrete Fourier transform to define $\sigma_A$.

In section 5 we consider operators in the Wiener algebra $W(X)$ with slowly oscillating coefficients. These operators are distinguished by two remarkable properties: their limit operators are periodic operators, and all limit operators belong to the Wiener algebra again. Via formula (1) we thus obtain a complete description of the essential spectra of operators with slowly oscillating coefficients.

In section 6 we apply these results to Schrödinger operators with slowly oscillating electrical potentials. As already mentioned, every $\mathbb{Z}^n$-periodic graph induces a related Schrödinger operator in a natural way (it is only this place where the graph structure becomes important). As illustrations we calculate the essential spectra of Schrödinger operators with slowly oscillating potentials on the zig-zag graph and on the hexagonal graph. The spectra of periodic Hamiltonians on zig-zag and hexagonal graphs connected with carbon nano-structures were considered in [13, 14, 19]. Note that slowly oscillating perturbations of periodic operators can change the spectrum of the unperturbed operator drastically. For instance, gaps in the spectrum of the unperturbed operator can be closed by the essential spectrum of the perturbed operator.

In section 7 we examine the essential spectrum of the Hamiltonian of the motion of two particles on a periodic graph $\Gamma$ around a heavy nucleus. For the lattice $\Gamma = \mathbb{Z}^n$ we considered this problem in [31]. See also the papers [1, 2, 21, 25, 26] and the references therein which are devoted to discrete multi-particle problems.

The limit operators approach does also apply to study the essential spectrum of pseudodifferential operators on periodic quantum graphs. We plan to develop these ideas in a forthcoming paper.

The authors are grateful for the support by CONACYT (Project 43432) and by the German Research Foundation (Grant 444 MEX-112/2/05).

2. Band-dominated operators on $\mathbb{Z}^n$

In this section we fix some notations and recall some facts concerning the Fredholm property of band-dominated operators on $l^p(\mathbb{Z}^n)$. The Fredholm properties of these operators are fairly well understood. All details can be found in [33]; see also the monograph [35] for a comprehensive account.

We will use the following notations. Given a Banach space $X$, let $\mathcal{L}(X)$ refer to the Banach algebra of all bounded linear operators on $X$ and $\mathcal{K}(X)$ to the closed ideal of the compact operators. An operator $A \in \mathcal{L}(X)$ is called a Fredholm operator if its kernel $\ker A := \{x \in X : Ax = 0\}$ and its cokernel $\text{coker} A := X/\mathcal{K}(X)$ are finite dimensional. Equivalently, $A$ is Fredholm if the coset $A + \mathcal{K}(X)$ is invertible in the Calkin algebra $\mathcal{L}(X)/\mathcal{K}(X)$. The essential spectrum of $A$ is the set of all complex numbers $\lambda$ for which the operator $A - \lambda I$ is not Fredholm on $X$, whereas the discrete spectrum of $A$ consists of all isolated eigenvalues of finite multiplicity. We denote the essential spectrum of $A$ by $\text{sp}_{\text{ess}} A$, the discrete spectrum by $\text{sp}_{\text{dis}} A$ and the usual spectrum by $\text{sp} A$. Sometimes we also write
sp(\(A : X \to X\)) instead of sp A in order to emphasize the underlying space X (with obvious modifications for the essential and the discrete spectrum). Clearly,

\[ \text{sp}_{\text{ess}}(A) \subseteq \text{sp}(A) \setminus \text{sp}_{\text{disc}}(A) \]

for every operator A \(\in \mathcal{L}(X)\). If A is self-adjoint, then equality holds in this inclusion.

Let \(p \geq 1\) be a real number and \(n\) a positive integer. As usual, we write \(l^p(\mathbb{Z}^n)\) for the Banach space of all functions \(u : \mathbb{Z}^n \to \mathbb{C}\) for which

\[ \|u\|_{l^p(\mathbb{Z}^n)} := \left( \sum_{x \in \mathbb{Z}^n} |u(x)|^p \right)^{1/p} < \infty \]

and \(l^\infty(\mathbb{Z}^N)\) for the Banach space of all bounded functions \(u : \mathbb{Z}^n \to \mathbb{C}\) with norm

\[ \|u\|_{l^\infty(\mathbb{Z}^n)} := \sup_{x \in \mathbb{Z}^n} |u(x)|. \]

For every positive integer \(N\), let \(l^p(\mathbb{Z}^n)^N\) stand for the Banach space of all vectors \(u = (u_1, \ldots, u_N)\) of functions \(u_i \in l^p(\mathbb{Z}^n)\) with norm

\[ \|u\|_{l^p(\mathbb{Z}^n)^N} := \left( \sum_{i=1}^N \|u_i\|_{l^p(\mathbb{Z}^n)}^p \right)^{1/p}. \]

Likewise, one can identify \(l^p(\mathbb{Z})^N\) with the Banach space \(l^p(\mathbb{Z}^n, \mathbb{C}^N)\) of all functions \(u : \mathbb{Z}^n \to \mathbb{C}^N\) for which

\[ \|u\|_{l^p(\mathbb{Z}^n, \mathbb{C}^N)} := \left( \sum_{x \in \mathbb{Z}^n} \sum_{i=1}^N |u_i(x)|^p \right)^{1/p} < \infty. \]

Clearly, the Banach spaces \(l^p(\mathbb{Z}^n)^N\) and \(l^p(\mathbb{Z}^n, \mathbb{C}^N)\) are isometric to each other. We also consider the Banach spaces \(l^\infty(\mathbb{Z}^n)^N\) and \(l^\infty(\mathbb{Z}^n, \mathbb{C}^N)\) with norms

\[ \|u\|_{l^\infty(\mathbb{Z}^n)^N} := \sup_{1 \leq i \leq N} \|u_i\|_{l^\infty(\mathbb{Z}^n)} \]

and

\[ \|u\|_{l^\infty(\mathbb{Z}^n, \mathbb{C}^N)} := \sup_{x \in \mathbb{Z}^n} \sup_{1 \leq i \leq N} |u_i(x)|. \]

Again, these spaces are isometric to each other in a natural way. Note also that \(l^\infty(\mathbb{Z}^n, \mathbb{C}^{N \times N})\) can be made to a C*-algebra by providing the matrix algebra \(\mathbb{C}^{N \times N}\) with a C*-norm.

We consider operators on \(l^p(\mathbb{Z}^n, \mathbb{C}^N)\) which are constituted by shift operators and by operators of multiplication by bounded functions. The latter are defined as follows: For \(\alpha \in \mathbb{Z}^n\), the shift operator \(V_\alpha\) is the isometry acting on \(l^p(\mathbb{Z}^n, \mathbb{C}^N)\) by \((V_\alpha u)(x) := u(x - \alpha)\). Further, each function \(a\) in \(l^\infty(\mathbb{Z}^n, \mathbb{C}^{N \times N})\) induces a multiplication operator \(a I\) on \(l^p(\mathbb{Z}^n, \mathbb{C}^N)\) via \((a I)(x) := a(x)u(x)\). Clearly,

\[ \|a I\|_{\mathcal{L}(l^p(\mathbb{Z}^n, \mathbb{C}^N))} = \|a\|_{l^\infty(\mathbb{Z}^n, \mathbb{C}^{N \times N})}. \]

A band operator on \(l^p(\mathbb{Z}^n, \mathbb{C}^N)\) is an operator of the form

\[ A = \sum_{|\alpha| \leq m} a_\alpha V_\alpha \]

with coefficients \(a_\alpha \in l^\infty(\mathbb{Z}^n, \mathbb{C}^{N \times N})\). The closure in \(\mathcal{L}(l^p(\mathbb{Z}^n, \mathbb{C}^N))\) of the set of all band operators is a subalgebra of \(\mathcal{L}(l^p(\mathbb{Z}^n, \mathbb{C}^N))\). We denote this algebra by \(A(l^p(\mathbb{Z}^n, \mathbb{C}^N))\) and call its elements band-dominated operators (BDO for short). Analogously, band-dominated operators on \(l^\infty(\mathbb{Z}^n, \mathbb{C}^N)\) are defined.
Our main tool to study the Fredholm property of band-dominated operators are the associated limit operators.

**Definition 1.** Let \( A \in \mathcal{L}(l^p(\mathbb{Z}^n, \mathbb{C}^N)) \), and let \( h : \mathbb{N} \to \mathbb{Z}^n \) be a sequence tending to infinity. A linear operator \( A_h \) is called the limit operator of \( A \) with respect to the sequence \( h \) if

\[
V^{-h(m)}AV_{h(m)} \to A_h \quad \text{and} \quad V^{-h(m)}A^*V_{h(m)} \to A_h^* \text{ strongly as } m \to \infty.
\]

We let \( \text{op}_p A \) denote the set of all limit operators of \( A \).

Here and in what follows, convergence of a sequence in \( \mathbb{Z}^n \) to infinity means convergence of this sequence to infinity in the one-point compactification of \( \mathbb{Z}^n \).

There are operators on \( l^p(\mathbb{Z}^n, \mathbb{C}^N) \) which do not possess limit operators at all. But if \( A \) is a band-dominated operator then one can show via a Cantor diagonal argument that every sequence \( h \) tending to infinity has a subsequence \( g \) for which the limit operator \( A_g \) exists. Moreover, the operator spectrum of \( A \) stores the complete information on the Fredholm property of \( A \), as the following theorem states. (In the case of \( n = 1 \) there is also a nice formula for the Fredholm index of \( A \); see [32].)

**Theorem 2.** An operator \( A \in \mathcal{A}(l^p(\mathbb{Z}^n, \mathbb{C}^N)) \) is Fredholm if and only if all limit operators of \( A \) are invertible and if

\[
\sup_{A_h \in \text{op}_p(A)} \| A_h^{-1} \| < \infty.
\] (3)

The uniform boundedness condition (3) is often difficult to check: It is one thing to verify the invertibility of an operator and another one to provide a good estimate for the norm of its inverse. It is therefore of vital importance to single out classes of band-dominated operators for which this condition is automatically satisfied. One of these classes is defined by imposing conditions on the decay of the norms of the coefficients. More precisely, we consider band-dominated operators of the form

\[
A := \sum_{a \in \mathbb{Z}^n} a_{\alpha} V_{\alpha},
\]

where

\[
\sum_{a \in \mathbb{Z}^n} \| a_{\alpha} \|_{l^\infty(\mathbb{Z}^n, \mathbb{C}^{n \times n})} < \infty.
\] (4)

One can show that the set \( W(\mathbb{Z}^n, \mathbb{C}^N) \) of all operators with property (4) forms an algebra and that the term on the left-hand side of (4) defines a norm which makes \( W(\mathbb{Z}^n, \mathbb{C}^N) \) to a Banach algebra. We refer to this algebra as the Wiener algebra and write \( \| A \|_{W(\mathbb{Z}^n, \mathbb{C}^N)} \) for the norm of an operator in \( W(\mathbb{Z}^n, \mathbb{C}^N) \). Clearly, operators in the Wiener algebra are bounded on each of the spaces \( l^p(\mathbb{Z}^n, \mathbb{C}^N) \) (including \( p = \infty \)) and

\[
\| A \|_{l^p(\mathbb{Z}^n, \mathbb{C}^N)} \leq \| A \|_{W(\mathbb{Z}^n, \mathbb{C}^N)}.
\]

Hence, \( W(\mathbb{Z}^n, \mathbb{C}^N) \subseteq \mathcal{A}(l^p(\mathbb{Z}^n, \mathbb{C}^N)) \) for every \( p \).

One important property of the Wiener algebra is its inverse closedness in each of the algebras \( \mathcal{L}(l^p(\mathbb{Z}^n, \mathbb{C}^N)) \), i.e., if \( A \in W(\mathbb{Z}^n, \mathbb{C}^N) \) has an inverse in \( \mathcal{L}(l^p(\mathbb{Z}^n, \mathbb{C}^N)) \), then \( A^{-1} \) belongs to \( W(\mathbb{Z}^n, \mathbb{C}^N) \) again. This fact implies that the spectrum of an operator \( A \in W(\mathbb{Z}^n, \mathbb{C}^N) \) considered as acting on \( l^p(\mathbb{Z}^n, \mathbb{C}^N) \) does not depend on \( p \in (1, \infty) \). Also the operator spectrum \( \text{op}_p(A) \) proves to be independent of \( p \), which justifies to write \( \text{op} A \).
instead. Note finally that all limit operators of operators in the Wiener algebra belong to the
Wiener algebra again.

For operators in the Wiener algebra, the Fredholm criterion in theorem 2 reduces to the
following much simpler assertion.

**Theorem 3.** Let $A \in W(\mathbb{Z}^n, \mathbb{C}^N)$. The operator $A$ is Fredholm on $L^p(\mathbb{Z}^n, \mathbb{C}^N)$ if and only if there exists a $p_0 \in [1, \infty]$ such that all limit operators of $A$ are invertible on $L^{p_0}(\mathbb{Z}^n, \mathbb{C}^N)$.

Theorem 3 has the following useful consequence.

**Theorem 4.** For $A \in W(\mathbb{Z}^n, \mathbb{C}^N)$, the essential spectra of $A : L^p(\mathbb{Z}^n, \mathbb{C}^N) \to L^p(\mathbb{Z}^n, \mathbb{C}^N)$ do not depend on $p \in (1, \infty)$, and

$$\text{sp}_{\text{ess}} A = \bigcup_{A_h \in \text{op } A} \text{sp } A_h.$$  \hfill (5)

3. BDO on periodic discrete metric spaces

3.1. Periodic discrete metric spaces

By a discrete metric space we mean a countable set $X$ together with a metric $\rho$ such that every ball

$$B_r(x_0) := \{ x \in X : \rho(x, x_0) \leq r \}$$

is a finite set. For each discrete metric space $X$, we introduce some standard Banach spaces over $X$. For $p \in (1, \infty)$, let $L^p(X)$ denote the Banach space of all complex-valued functions $u$ on $X$ with norm

$$\|u\|_{L^p(X)} := \sum_{x \in X} |u(x)|^p,$$

and write $L^\infty(X)$ for the Banach space of all bounded functions $u$ of $X$ with norm

$$\|u\|_{L^\infty(X)} := \sup_{x \in X} |u(x)|.$$

A periodic discrete metric space is a discrete metric space provided with the free action of the group $\mathbb{Z}^n$. More precisely, let $X$ be a discrete metric space, and let there be a mapping

$$\mathbb{Z}^n \times X \to X, \quad (\alpha, x) \mapsto \alpha \cdot x$$

satisfying

$$0 \cdot x = x \quad \text{and} \quad (\alpha + \beta) \cdot x = \alpha \cdot (\beta \cdot x)$$

for arbitrary elements $\alpha, \beta \in \mathbb{Z}^n$ and $x \in X$, which leaves the metric invariant,

$$\rho(\alpha \cdot x, \alpha \cdot y) = \rho(x, y)$$ \hfill (6)

for all elements $\alpha \in \mathbb{Z}^n$ and $x, y \in X$. Recall also that the group $\mathbb{Z}^n$ acts freely on $X$ if whenever the equality $x = \alpha \cdot x$ holds for elements $x \in X$ and $\alpha \in \mathbb{Z}^n$ then, necessarily, $\alpha = 0$.

For each element $x \in X$, consider its orbit $\{ \alpha \cdot x : \alpha \in \mathbb{Z}^n \}$ with respect to the action of $\mathbb{Z}^n$. Any two orbits are either disjoint or identical. Hence, there is a binary equivalence relation on $X$, by calling two points equivalent if they belong to the same orbit. The set of all orbits of $X$ with respect to the action of $\mathbb{Z}^n$ is denoted by $X/\mathbb{Z}^n$. A basic assumption
throughout what follows is that the orbit space $X/\mathbb{Z}^n$ is finite. Thus, there is a finite subset $M := \{x_1, x_2, \ldots, x_N\}$ of $X$ such that the orbits

$$X_j := \{\alpha \cdot x_j \in X : \alpha \in \mathbb{Z}^n\}$$

satisfy $X_i \cap X_j = \emptyset$ if $x_i \neq x_j$ and $\bigcup_{i=1}^N X_i = X$. If all these conditions are satisfied then we call $X$ a periodic discrete metric space with respect to $\mathbb{Z}^n$ or simply $\mathbb{Z}^n$-periodic.

The free action of $\mathbb{Z}^n$ on $X$ guarantees that the mapping

$$U_j : \mathbb{Z}^n \rightarrow X_j, \quad \alpha \mapsto \alpha \cdot x_j$$

is a bijection for every $j = 1, \ldots, N$. For each complex-valued function $f$ on $X$, let $Uf : \mathbb{Z}^n \rightarrow \mathbb{C}^N$ be the function

$$(Uf)(\alpha) := ((U_1 f)(\alpha), \ldots, (U_N f)(\alpha)).$$

Clearly, the mapping $U$ is a linear isometry from $l^p(X)$ onto $l^p(\mathbb{Z}^n, \mathbb{C}^N)$, and the mapping $A \mapsto UA^{-1}$ is an isometric isomorphism from $L(l^p(X))$ onto $L(l^p(\mathbb{Z}^n, \mathbb{C}^N))$ for every $p \in [1, \infty]$.

Another consequence of our assumptions is that

$$\lim_{\mathbb{Z}^n \ni \alpha \rightarrow \infty} \rho(\alpha \cdot x, y) = \infty$$

for all points $x, y \in X$. Indeed, suppose that (7) is wrong. Then there are points $x, y \in X$, a positive constant $M$, and a sequence $\alpha$ of pairwise different points in $\mathbb{Z}^n$ such that

$$\rho(\alpha(n) \cdot x, y) \leq M \quad \text{for all} \ n \in \mathbb{N}. \quad (8)$$

The free action of $\mathbb{Z}^n$ on $X$ implies that $(\alpha(n) \cdot x)_{n \in \mathbb{N}}$ is a sequence of pairwise different points in $X$. Hence, (8) implies that the ball with center $y$ and radius $M$ contains infinitely many points, a contradiction.

### 3.2. Band-dominated operators on $X$

Let $X$ be a periodic discrete metric space and $p \in [1, \infty)$. We consider linear operators $A$ on $l^p(X)$ for which there exists a function $k_A \in l^\infty(X \times X)$ such that

$$(Au)(x) = \sum_{y \in X} k_A(x, y)u(y) \quad \text{for all} \ x \in X \quad (9)$$

and for all finitely supported functions $u$ on $X$ (note that the latter form a dense subspace of $l^p(X)$). We call $k_A$ the generating function of the operator $A$. It is easily seen that every bounded operator $A$ on $l^p(X)$ is of this form and is, thus, generated by a bounded function. The converse is certainly not true. It is also clear that every operator $A$ determines its generating function uniquely, since

$$(A\delta_y)(x) = k_A(x, y),$$

where $\delta_y$ is the function on $X$ which is 1 at $y$ and 0 at all other points.

An operator $A$ of the form (9) is called a band operator if there exists an $R > 0$ such that $k_A(x, y) = 0$ whenever $\rho(x, y) > R$.

**Example 5.** Every operator $aI$ of multiplication by a function $a \in l^\infty(X)$ is a band operator.

**Example 6.** For $\alpha \in \mathbb{Z}^n$, let $T_{\alpha}$ be the operator of shift by $\alpha$ on $l^p(X)$, i.e., $(T_{\alpha}u)(x) := u(\cdot - \alpha) \cdot x)$. Clearly, $T_{\alpha}$ is a band operator which acts as an isometry on $l^p(X)$. Hence, every operator of the form

$$\sum_{|\alpha| \leq m} a_{\alpha} T_{\alpha}$$

(10)
with $a_u \in l^\infty(X)$ is a band operator (but there are band operators which cannot be represented of this form).

**Proposition 7.** If $A$ is a band operator on $l^p(X)$, then $UAU^{-1}$ is a band operator on $l^p(\mathbb{Z}^n, \mathbb{C}^N)$.

**Proof.** The operator $UAU^{-1}$ has the matrix representation

$$
(UAU^{-1} f)(\alpha) = \sum_{j=1}^N \sum_{\beta \in \mathbb{Z}^n} r^j_\alpha (\alpha, \beta) f_j (\beta),
$$

(11)

where $\alpha \in \mathbb{Z}^n, i = 1, \ldots, N$ and

$$
r^j_\alpha (\alpha, \beta) := k_A(\alpha \cdot x_i, \beta \cdot x_j).
$$

(12)

From (7) we conclude that

$$
\rho(\alpha \cdot x_i, \beta \cdot x_j) = \rho(x_i, (\beta - \alpha) \cdot x_j) \to \infty
$$
as $|\alpha - \beta| \to \infty$. Thus, there is an $R_1 > 0$ such that $r^j_\alpha (\alpha, \beta) = 0$ if $|\alpha - \beta| > R_1$. In other words, every $r^j_\alpha$ is the generating function of a band operator on $l^p(\mathbb{Z}^n)$, implying that $UAU^{-1}$ is a band operator on $l^p(\mathbb{Z}^n, \mathbb{C}^N)$. □

The preceding proposition implies in particular that every band operator is bounded on $l^p(X)$ for every $p \in [1, \infty]$. Let $A_p(X)$ stand for the closure in $L(l^p(X))$ of the set of all band operators. The operators in $A_p(X)$ are called band-dominated operators on $X$. Note that $A_p(X)$ depends heavily on $p$ whereas the class of the band operators is independent of $p$. One can show easily (for example, by employing the preceding proposition and the well-known properties of band-dominated operators on $\mathbb{Z}^n$) that $A_p(X)$ is a Banach algebra for every $p$ and a $C^*$-algebra for $p = 2$.

**Proposition 8.** Let $X$ be a periodic discrete metric space and $p \in [1, \infty]$. The mapping $A \mapsto UAU^{-1}$ is an isomorphism between the Banach algebras $A_p(X)$ and $A_p(\mathbb{Z}^n, \mathbb{C}^N)$.

**Proof.** Note that an operator $A$ is a band operator on $l^p(X)$ if and only if $UAU^{-1}$ is a band operator on $l^p(\mathbb{Z}^n)$, implying that $UAU^{-1}$ is a band operator on $l^p(\mathbb{Z}^n, \mathbb{C}^N)$. Thus, the assertion follows since the mapping $A \mapsto UAU^{-1}$ is a continuous isomorphism between the Banach algebras $L(l^p(X))$ and $L(l^p(\mathbb{Z}^n, \mathbb{C}^N))$. □

### 3.3. Limit operators and Fredholm property

Let $X$ be a $\mathbb{Z}^n$-periodic discrete metric space. The goal of this section is a criterion for the Fredholm property of band-dominated operators on $l^p(X)$. This criterion makes use of the limit operators of $A$ which, in a sense, reflect the behavior of $A$ at infinity. Here is the definition.

**Definition 9.** Let $1 < p < \infty$, and $h : \mathbb{N} \to \mathbb{Z}^n$ be a sequence tending to infinity. We say that $A_h$ is a limit operator of $A \in L(l^p(X))$ defined by the sequence $h$ if

$$
T_{h(m)}^{-1} A T_{h(m)} \to A_h \quad \text{and} \quad T_{h(m)}^{-1} A^* T_{h(m)} \to A_h^* \quad \text{as} \quad m \to \infty
$$

strongly on $l^p(X)$ and $l^q(X)^* = l^q(X)$ with $1/p + 1/q = 1$, respectively. We denote the set of all limit operators of $A$ by $\text{op}_p(A)$ and call this set the operator spectrum of $A$.

Note that the generating function of the shifted operator $T_u^{-1} A T_u$ is related with that of $A$ by

$$
k_{T_u^{-1} A T_u}(x, y) = k_A((-\alpha) \cdot x, (-\alpha) \cdot y)
$$

(13)
and that the generating functions of $T_{h(m)}^{-1}A_{h(m)}$ converge point-wise on $X \times X$ to the generating function of the limit operator $A_h$ if the latter exists.

It is an important property of band-dominated operators that their operator spectrum is not empty. More general, one has the following result which can be proved by an obvious Cantor diagonal argument (see [33–35]).

**Proposition 10.** Let $p \in (1, \infty)$ and $A \in \mathcal{A}_p(X)$. Then every sequence $h : \mathbb{N} \to G$ which tends to infinity possesses a subsequence $g$ such that the limit operator $A_g$ of $A$ with respect to $g$ exists.

The following theorem settles the basic relation between the Fredholm property of a band-dominated operator $A$ and the invertibility of its limit operators. It follows easily from theorem 2 if one takes into account that the mapping

$$\mathcal{A}_p(X) \to \mathcal{A}_p(\mathbb{Z}^n, \mathbb{C}^N), \quad A \mapsto UAU^{-1}$$

is an isomorphism of Banach algebras and that the relation

$$(UAU^{-1})_h = U A_h U^{-1}$$

between the limit operators of $A$ and $UAU^{-1}$ holds.

**Theorem 11.** Let $p \in (1, \infty)$ and $A \in \mathcal{A}_p(X)$. Then $A$ is a Fredholm operator on $l^p(X)$ if and only if all limit operators of $A$ are invertible and if the norms of their inverses are uniformly bounded,

$$\sup_{A_h \in \text{op}(A)} \|A_h^{-1}\| < \infty. \quad (14)$$

### 3.4. The Wiener algebra of $X$

The goal of this section is to single out a class of band-dominated operators for which the uniform boundedness condition (14) is redundant.

**Definition 12.** Let $X$ be a $\mathbb{Z}^n$-periodic discrete metric space. The set $\mathcal{W}(X)$ consists of all linear operators $A$ for which there is a function $h_A$ in $l^1(\mathbb{Z}^n)$ such that

$$\max_{j \in \{1, \ldots, N\}} \sum_{i=1}^N |r_{ij}^1(\alpha, \beta)| \leq h_A(\alpha - \beta) \quad (15)$$

for all $\alpha, \beta \in \mathbb{Z}^n$.

We introduce a norm in $\mathcal{W}(X)$ by

$$\|A\|_{\mathcal{W}(X)} := \inf \|h\|_{l^1(\mathbb{Z}^n)}, \quad (16)$$

where the infimum is taken over all sequences $h \in l^1(\mathbb{Z}^n)$ for which inequality (15) holds in place of $h_A$.

**Proposition 13.** The set $\mathcal{W}(X)$ with the norm (16) is a Banach algebra, and the mapping $A \mapsto UAU^{-1}$ is an isometric isomorphism between the Banach algebras $\mathcal{W}(X)$ and $\mathcal{W}(\mathbb{Z}^n, \mathbb{C}^N)$.

The proof is straightforward. We refer to the algebra $\mathcal{W}(X)$ as the Wiener algebra.

**Proposition 14.** Let $p \in [1, \infty]$.

(i) Every operator $A \in \mathcal{W}(X)$ is bounded on each of the spaces $l^p(X)$.

(ii) The algebra $\mathcal{W}(X)$ is inverse closed in each of the algebras $\mathcal{L}(l^p(X))$. 
Proposition 14 follows from proposition 13 and the related results for the special case $X = \mathbb{Z}^n$ presented in [33, 34] and [35].

The following result highlights the importance of the Wiener algebra in our context.

**Theorem 15.** Let $A \in \mathcal{W}(X)$. Then $A$ is a Fredholm operator on $l^p(X)$ with $p \in (1, \infty)$ if and only if there is a $p_0 \in [1, \infty]$ such that all limit operators of $A$ are invertible on $l^{p_0}(X)$. Moreover, $\text{spess } A$ does not depend on $p \in (1, \infty)$, and

$$\text{spess } A = \bigcup_{A_h \in \text{cop}(A)} \text{sp } A_h. \quad (17)$$

Theorem 15 follows immediately from proposition 13 and theorems 3 and 4.

**Corollary 16.** Let $B := A + aI$ where $A \in \mathcal{W}(X)$ and $\lim_{x \to \infty} a(x) = 0$. Then $\text{spess } B = \text{spess } A$.

The following result states a sufficient condition for the absence of the discrete spectrum of an operator $A \in \mathcal{A}_p(X)$.

**Proposition 17.** Let $A \in \mathcal{A}_p(X)$ and suppose there is a sequence $h : \mathbb{N} \to \mathbb{Z}^n$ for which the limit operator $A_h$ exists in the sense of norm convergence,

$$\lim_{m \to \infty} \|T_h^{-1} A T_h - A_h\| = 0. \quad (18)$$

Then $\text{spess } A = \text{sp } A$.

**Proof.** Let $\lambda \notin \text{spess } A$. Then, by theorem 11, $\lambda \notin \text{sp } A_h$. It follows from (18) that $\lambda \notin \text{sp } A$. Hence, $\text{sp } A \subseteq \text{spess } A$, which implies the assertion. \qed

### 4. Periodic operators on periodic metric spaces

Let $X$ be a $\mathbb{Z}^n$-periodic discrete metric space. An operator $A \in \mathcal{L}(l^p(X))$ is $\mathbb{Z}^n$-periodic if it is invariant with respect to left shifts by elements of $\mathbb{Z}^n$,

$$T_\alpha A = A T_\alpha \quad \text{for every } \alpha \in \mathbb{Z}^n.$$

The following is a straightforward consequence of proposition 17.

**Proposition 18.** Let $A \in \mathcal{A}_p(X)$ be a $\mathbb{Z}^n$-periodic operator. Then $\text{spess } A = \text{sp } A$.

Similar results are well known for periodic differential operators (see, e.g. [15]).

An explicit description of the spectrum of $\mathbb{Z}^n$-periodic operators can be given via the Fourier transform. One easily checks that $A \in \mathcal{W}(X)$ is $\mathbb{Z}^n$-periodic if and only if the generating function $k_A$ of $A$ is periodic in the sense that, for all $\gamma \in \mathbb{Z}^n$ and all points $x, y \in X$,

$$k_A(\gamma \cdot x, \gamma \cdot y) = k_A(x, y).$$

Periodicity thus implies that the functions $r_A^j (\alpha, \beta) := k_A(\alpha \cdot x_j, \beta \cdot x_j)$ satisfy

$$r_A^j (\alpha, \beta) = k_A((\alpha - \gamma) \cdot x_i, (\beta - \gamma) \cdot x_j)$$
Difference operators on periodic graphs

for all \( y \in \mathbb{Z}^n \), whence \( r^j_i(\alpha, \beta) = r^j_i(\alpha - \beta, 0) \). Hence, for \( i = 1, \ldots, N \),

\[
(UA^{-1}f)_i(\alpha) = \sum_{j=1}^{N} \sum_{\beta \in \mathbb{Z}^n} r^j_i(\alpha, \beta)(U_j f)(\beta)
\]

\[
= \sum_{j=1}^{N} \sum_{\beta \in \mathbb{Z}^n} r^j_i(\beta, 0)(V_{\beta}U_j f)(\alpha),
\]

where \( |r^j_i(\beta, 0)| \leq h(\beta) \) for some non-negative function \( h \in l^1(\mathbb{Z}^n) \). Thus, we obtained the following.

**Proposition 19.** Every \( \mathbb{Z}^n \)-periodic operator \( A \in \mathcal{W}(X) \) is isometrically equivalent to the shift invariant matrix operator \( UA^{-1} \in \mathcal{W}(\mathbb{Z}^n, \mathbb{C}^N) \).

Under the conditions of the previous proposition, we associate with \( A \) a function \( \sigma_A : T^m \rightarrow \mathbb{C}^N \times N \), called the symbol of \( A \), via

\[
\sigma_A(t) := \sum_{\beta \in \mathbb{Z}^n} r_A(\beta)t^\beta,
\]

where \( T \) is the torus \( \{z \in \mathbb{C} : |z| = 1\} \), \( r_A(\beta) \) is the matrix \( \left(r^j_i(\beta, 0)\right)_{j=1}^{N} \), and \( t^\beta := t_1^{\beta_1} \cdots t_n^{\beta_n} \) for \( t = (t_1, \ldots, t_n) \in T^m \) and \( \beta = (\beta_1, \ldots, \beta_n) \in \mathbb{Z}^n \). It is well known that the operator

\[
(\tilde{A}u)(\alpha) := \sum_{\beta \in \mathbb{Z}^n} r_A(\alpha - \beta, 0)u(\beta)
\]

is invertible on \( l^p(\mathbb{Z}^n, \mathbb{C}^N) \) with \( p \in [1, \infty] \) if and only if det \( \sigma_A \neq 0 \) on \( T^m \).

For \( t \in T^m \), let \( \lambda^j_A(t) \) with \( j = 1, \ldots, N \) denote the eigenvalues of the matrix \( \sigma_A(t) \). The enumeration of the eigenvalues can be chosen in such a way that \( \lambda^j_A(t) \) depends continuously on \( t \) for every \( j \). The sets

\[
C_j(A) := \{ \lambda \in \mathbb{C} : \lambda = \lambda^j_A(t), t \in T^m \}, \quad j = 1, \ldots, N
\]

are called the spectral or dispersion curves of \( A \).

**Proposition 20.** Let \( A \in \mathcal{W}(X) \) be a \( \mathbb{Z}^n \)-periodic operator. Then

\[
\text{sp } A = \text{sp ess } A = \bigcup_{j=1}^{N} C_j(A).
\]

### 5. Operators with slowly oscillating coefficients on periodic metric spaces

Let again \( X \) be a \( \mathbb{Z}^n \)-periodic discrete metric space. A function \( a \in l^\infty(X) \) is slow oscillating if, for every two points \( x, y \in X \),

\[
\lim_{a \to \infty} (a(\alpha \cdot x) - a(\alpha \cdot y)) = 0.
\]

The set of all slowly oscillating functions on \( X \) forms a \( C^* \)-subalgebra of \( l^\infty(X) \) which we denote by \( \text{SO}(X) \). Note that the class \( \text{SO}(X) \) does not only depend on \( X \) but also on the action of \( \mathbb{Z}^n \) on \( X \).

Let \( a \in \text{SO}(X) \) and \( h : \mathbb{N} \rightarrow G \) be a sequence tending to infinity. The Bolzano–Weierstrass theorem and a Cantor diagonal argument imply that there is a subsequence \( g \) of \( h \) such that the functions \( x \mapsto a(g(m) \cdot x) \) converge point-wise to a function \( a_g \in l^\infty(X) \) as
We consider the operators of the form
\[ A = \sum_{k,l=1}^{\infty} b_k A_{kl} c_l I, \] (22)
where the \( A_{kl} \) are \( \mathbb{Z}^n \)-periodic operators in \( \mathcal{W}(X) \) and the \( b_k \) and \( c_l \) are slowly oscillating functions satisfying
\[ \sum_{k,l=1}^{\infty} \| b_k \|_{L^\infty(X)} \| A_{kl} \|_{\mathcal{W}(X)} \| c_l \|_{L^\infty(X)} < \infty. \]

Let \( h : \mathbb{N} \to \mathbb{Z}^n \) be a sequence tending to infinity. Then
\[ T_{h(m)}^{-1} A T_{h(m)} = \sum_{k,l=1}^{\infty} (T_{h(m)}^{-1} b_k) A_{kl} (T_{h(m)}^{-1} c_l) I. \]

One can assume without loss that the point-wise limits
\[ \lim_{m \to \infty} (T_{h(m)}^{-1} b_k)(x) =: b_k^h, \quad \lim_{m \to \infty} (T_{h(m)}^{-1} c_l)(x) =: c_l^h \]
exist (otherwise we pass to a suitable subsequence of \( h \)). As we have seen above, the limit functions \( b_k^h \) and \( c_l^h \) are \( \mathbb{Z}^n \)-periodic on \( X \). Consequently, the limit operators \( A_h \) of \( A \) are \( \mathbb{Z}^n \)-periodic operators of the form
\[ A_h = \sum_{k,l=1}^{\infty} b_k^h A_{kl} c_l^h I. \]

Now, the following is an immediate consequence of theorem 15.

**Theorem 21.** Let \( A \) be an operator with slowly oscillating coefficients of the form (22). Then \( A \) is a Fredholm operator on \( l^p(X) \) if and only if, for every operator \( A_h \in \text{op} A \),
\[ \det \sigma_{A_h}(t) \neq 0 \quad \text{for every} \quad t \in \mathbb{T}^n. \]

Moreover,
\[ \text{sp}_{\text{ess}} A = \bigcup_{A_h \in \text{op} A} \text{sp} A_h = \bigcup_{A_h \in \text{op} A} \bigcup_{j=1}^{N} C_j(A_h). \]

6. **Schrödinger operators on periodic graphs**

By a *discrete infinite graph* we mean a countable set \( X \) together with a binary relation \( \sim \) which is anti-reflexive (i.e., there is no \( x \in X \) such that \( x \sim x \)) and symmetric and which has the property that for each \( x \in X \) there are only finitely many \( y \in X \) such that \( x \sim y \). The points of \( X \) are called the *vertices* and the pairs \((x, y)\) with \( x \sim y \) the *edges* of the graph. Due to anti-reflexivity, the graphs under consideration do not possess loops. We write \( m(x) \) for the number of edges starting (or ending) at the vertex \( x \) of \( X \). If \( x \sim y \), we say that the vertices \( x, y \) are *adjacent*.

For technical reasons it will be convenient to assume that the graph \((X, \sim)\) is connected, i.e., given distinct points \( x, y \in X \), there are finitely many points \( x_0, x_1, \ldots, x_n \in X \) such that...
for arbitrary vertices $x, y \in X$ and group elements $\alpha \in \mathbb{Z}^n$. Clearly, every group with these properties leaves the graph distance invariant, that is, $X$ becomes a $\mathbb{Z}^n$-periodic discrete metric space. If $(X, \sim)$ is a $\mathbb{Z}^n$-periodic graph, then the function $m$ is $\mathbb{Z}^n$-periodic, too, that is, $m(\alpha \cdot x) = m(x)$ for every $x \in X$ and $\alpha \in \mathbb{Z}^n$.

Every $\mathbb{Z}^n$-periodic discrete graph $\Gamma := (X, \sim)$ induces a canonical difference operator $\Delta_\Gamma$ on $L^p(X)$, called the (discrete) Laplace operator or Laplacian of $\Gamma$, via

$$
(\Delta_\Gamma u)(x) := \frac{1}{m(x)} \sum_{y = x} u(y), \quad x \in X.
$$

(23)

Evidently, $\Delta_\Gamma$ is a $\mathbb{Z}^n$-periodic band operator.

Let $v \in L^\infty(X)$. The operator $\mathcal{H}_\Gamma := \Delta_\Gamma + vI$ is referred to as the (discrete) Schrödinger operator with electric potential $v$ on the graph $X$. Given a sequence $h : \mathbb{N} \to \mathbb{Z}^n$ tending to infinity, there exist a subsequence $g$ of $h$ and a function $v^g \in L^\infty(X)$ such that $v(g(m) \cdot x) \to v^g(x)$ as $m \to \infty$ for every $x \in X$. It turns out that the operator

$$
\mathcal{H}_\Gamma^g := \Delta_\Gamma + v^g I
$$

is the limit operator of $\mathcal{H}_\Gamma$ defined by the sequence $g$ and that every limit operator of $\mathcal{H}_\Gamma$ is of this form. Thus, theorem 15 implies the following.

**Theorem 22.** The Schrödinger operator $\mathcal{H}_\Gamma := \Delta_\Gamma + vI$ with bounded potential $v$ is a Fredholm operator on $L^p(X)$ with $p \in (1, \infty)$ if and only if there is a $p_0 \in [1, \infty]$ such that all limit operators of $\mathcal{H}_\Gamma$ are invertible on $L^{p_0}(X)$. The essential spectrum of $\mathcal{H}_\Gamma$ does not depend on $p \in (1, \infty)$, and

$$
\text{sp}_{\text{ess}} \mathcal{H}_\Gamma = \bigcup_{\mathcal{H}_\Gamma^g \in \text{sp} \mathcal{H}_\Gamma} \text{sp} \mathcal{H}_\Gamma^g.
$$

(24)

For an explicit description of the essential spectrum of the Schrödinger operator $\mathcal{H}_\Gamma$, we first assume that $v$ is a periodic potential. Then $UvU^{-1}$ is the operator of multiplication by the diagonal matrix $\text{diag}(v(x_1), \ldots, v(x_N))$. Hence,

$$
U\mathcal{H}_\Gamma U^{-1} = \sum_{\alpha \in \{1, 0, 1\}^n} a_\alpha V_\alpha + \text{diag}(v(x_1), \ldots, v(x_N)),
$$

where the $a_\alpha$ are certain constant $N \times N$ matrices which depend on the structure of the graph $\Gamma$. Consequently,

$$
\sigma_{\mathcal{H}_\Gamma}(t) = \sum_{\alpha \in \{1, 0, 1\}^n} a_\alpha t^\alpha + \text{diag}(v(x_1), \ldots, v(x_N)), \quad t \in \mathbb{R}^n.
$$

If the potential $v$ is real valued, then $\mathcal{H}_\Gamma$ is self-adjoint on $l^2(X)$ and $\sigma_{\mathcal{H}_\Gamma}$ is a Hermitian function. From proposition 20 we conclude that

$$
\text{sp} \mathcal{H}_\Gamma = \bigcup_{j=1}^N C_j(\mathcal{H}_\Gamma),
$$

where $C_j(\mathcal{H}_\Gamma)$ is the real interval $[a_j, b_j]$ with $a_j := \min_{t \in \mathbb{R}^n} \lambda_j^\Gamma (t)$ and $b_j := \max_{t \in \mathbb{R}^n} \lambda_j^\Gamma (t)$.
Next we consider Schrödinger operators \( H_{\Gamma} = \Delta_{\Gamma} + v I \) with slowly oscillating potential \( v \). As we have seen in the previous section, all limit operators of \( H_{\Gamma} \) are of the form \( H_{\Gamma}^{\pm} = \Delta_{\Gamma} + v^\pm I \) with periodic potentials \( v^\pm \). Theorems 22 and 15 imply the following.

**Theorem 23.** Let \( H_{\Gamma} = \Delta_{\Gamma} + v I \) with \( v \in SO(X) \). Then
\[
\text{spess } H_{\Gamma} = \bigcup_{H_{\Gamma}^{\pm} \in \text{op}(H_{\Gamma})} \bigcup_{j=1}^{N} C_j(H_{\Gamma}^{\pm})
\]
with the spectral curves \( C_j(H_{\Gamma}^{\pm}) \) defined as in (19).

If the slowly oscillating potential \( v \) is real valued, then the spectral curves \( C_j(H_{\Gamma}^{\pm}) \) are (possibly overlapping) real intervals.

The following examples clarify the structure of the essential spectrum of Schrödinger operators on some special periodic graphs. The graphs under consideration are embedded into \( \mathbb{R}^n \) for some \( n \). This embedding allows one to consider the vertices of the graph as vectors and to use the linear structure of \( \mathbb{R}^n \) in order to describe the group action. Note that we consider slowly oscillating potentials in these examples. The case of periodic potentials was studied in [13, 14, 19].

**Example 24 (The zig-zag graph).** Let \( \Gamma = (X, \sim) \) be the zig-zag graph in the plane \( \mathbb{R}^2 \) as shown in figure 1. The graph \( \Gamma \) is periodic with respect to the action \( g \cdot x_n := x_{n+2g} \) of the group \( \mathbb{Z} \), and the set \( M = \{x_1, x_2\} \) of vertices represents the fundamental domain.

One should mention that, as a graph, the zig-zag graph is isomorphic to the Cayley graph of the group \( \mathbb{Z} \) and, in both cases, it is the same group \( \mathbb{Z} \) which acts on the graph. The difference lies in the way in which \( \mathbb{Z} \) acts. For the Cayley graph, the group element \( \alpha \) maps the vertex \( x \) to \( \alpha + x \), whereas \( \alpha \) maps \( x \) to \( 2\alpha + x \) for the zig-zag graph. The latter action is visualized by the zig-zag form.

The operator \( U \Delta_{\Gamma} U^{-1} \) has the matrix representation
\[
U \Delta_{\Gamma} U^{-1} = \frac{1}{2} \begin{pmatrix} 0 & I + V_{(1,0)} \\ I + V_{(-1,0)} & 0 \end{pmatrix}
\]
in the basis induced by \( M \). Hence,
\[
\sigma_{\Delta_{\Gamma}}(t) = \frac{1}{2} \begin{pmatrix} 0 & 1 + t \\ 1 + t^{-1} & 0 \end{pmatrix}, \quad t \in \mathbb{T},
\]
and a straightforward calculation shows that the spectral curves of \( \Delta_{\Gamma} \) are
\[
\{ \lambda \in \mathbb{C} : \lambda = \pm \cos^2 \varphi / 2, \varphi \in [0, 2\pi] \}.
\]
Hence, the spectrum of the Laplacian \( \Delta_{\Gamma} \) of the zig-zag graph is the interval \([-1, 1]\).
Next consider the Schrödinger operator \( \mathcal{H}_\Gamma := \Delta_\Gamma + vI \) with \( \mathbb{Z} \)-periodic potential \( v \). Thus, \( v \) is completely determined by its values on \( \mathcal{M} \), and we write \( v_1 := v(x_1) \) and \( v_2 := v(x_2) \). Then

\[
\sigma_{\mathcal{H}_\Gamma - \lambda I}(t) = \begin{pmatrix} v_1 - \lambda & (1 + t)/2 \\ (1 + t^{-1})/2 & v_2 - \lambda \end{pmatrix}, \quad t \in \mathbb{T},
\]

which implies that the spectral curves of \( \mathcal{H}_\Gamma \) are

\[
\lambda \in \mathbb{C} : \lambda = \frac{1}{2} \pm \frac{\sqrt{(v_1 - v_2)^2 + 4\cos^2 \varphi/2}}{2(v_1 + v_2)}, \quad \varphi \in [0, 2\pi].
\]

If, for example, \( v_1 \) and \( v_2 \) are real numbers with \( v_1 < v_2 \), then \( \text{sp}_{\text{ess}} \mathcal{H}_\Gamma = \text{sp} \mathcal{H}_\Gamma \) is the union of the disjoint intervals

\[
\left[ \frac{1}{2} - \frac{\sqrt{(v_1 - v_2)^2 + 4}}{2(v_1 + v_2)}, \frac{v_1}{v_1 + v_2} \right] \cup \left[ \frac{v_2}{v_1 + v_2}, \frac{1}{2} + \frac{\sqrt{(v_1 - v_2)^2 + 4}}{2(v_1 + v_2)} \right], \quad (25)
\]

that is, one observes a gap \( \left( \frac{v_1}{v_1 + v_2}, \frac{v_2}{v_1 + v_2} \right) \) in the spectrum.

Finally, let the potential \( v \) be slowly oscillating. Then the essential spectrum of \( \mathcal{H}_\Gamma \) is the union

\[
\bigcup_h \left[ \frac{1}{2} - \frac{\sqrt{(v^h_1 - v^h_2)^2 + 4}}{2(v^h_1 + v^h_2)}, \frac{\min \{v^h_1, v^h_2\}}{v^h_1 + v^h_2} \right] \cup \left[ \frac{\max \{v^h_1, v^h_2\}}{v^h_1 + v^h_2}, \frac{1}{2} + \frac{\sqrt{(v^h_1 - v^h_2)^2 + 4}}{2(v^h_1 + v^h_2)} \right], \quad (26)
\]

where the unions are taken with respect to all sequences \( h \) for which the limits

\[
v^h_j := \lim_{m \to \infty} v(h(m) \cdot x_j), \quad j = 1, 2, \quad (27)
\]

exist. Set

\[
a_{\mathcal{H}_\Gamma} := \limsup_{\mathbb{Z} \ni \alpha \to \infty} \frac{v(\alpha \cdot x_1)}{v(\alpha \cdot x_1) + v(\alpha \cdot x_2)},
\]

\[
b_{\mathcal{H}_\Gamma} := \liminf_{\mathbb{Z} \ni \alpha \to \infty} \frac{v(\alpha \cdot x_1) + v(\alpha \cdot x_2)}{v(\alpha \cdot x_1)}. \quad (28)
\]

Thus, if the inequality

\[
a_{\mathcal{H}_\Gamma} < b_{\mathcal{H}_\Gamma}
\]

holds, then the operator \( \mathcal{H}_\Gamma \) has the gap \( (a_{\mathcal{H}_\Gamma}, b_{\mathcal{H}_\Gamma}) \) in its essential spectrum. Of course, this interval can contain points of the discrete spectrum of \( \mathcal{H}_\Gamma \).

**Example 25.** [The honeycomb graph] Let \( \Gamma = (X, \sim) \) be the hexagonal graph shown in figure 2. We consider this graph as embedded into \( \mathbb{R}^2 \) and let \( e_1 \) and \( e_2 \) be the vectors indicated in the figure. The group \( \mathbb{Z}^2 \) operates on \( \Gamma \) via

\[
(\alpha_1, \alpha_2) \cdot x := x + \alpha_1 e_1 + \alpha_2 e_2
\]

(where \( \alpha_1, \alpha_2 \in \mathbb{Z} \) and \( x \in \mathbb{R} \)). A fundamental domain \( \mathcal{M} \) for this action is provided by any two vertices \( x_1, x_2 \) as marked in the figure.

Hence, we have to identify \( l^p(X) \) with \( l^p(\mathbb{Z}^2, \mathbb{C}^2) \), and the Laplacian \( \Delta_\Gamma \) has the following matrix representation with respect to \( \mathcal{M} \)

\[
U \Delta_\Gamma U^{-1} = \frac{1}{3} \begin{pmatrix} 0 & I + V e_1 + V e_2 \\ I + V e_1^{-1} + V e_2^{-1} & 0 \end{pmatrix}. \quad (29)
\]
Consequently,
\[ \sigma_{\Delta_G}(t) = \frac{1}{3} \left( \begin{array}{ccc} 0 & 1 + t_1 + t_2 & 1 + t_1^{-1} + t_2^{-1} \\ 1 + t_1^{-1} + t_2^{-1} & 0 & 1 + t_1 + t_2 \\ 1 + t_1 + t_2 & 1 + t_1^{-1} + t_2^{-1} & 0 \end{array} \right), \quad t = (t_1, t_2) \in \mathbb{T}^2, \]
and the spectral curves of the Laplacian \( \Delta_G \) are
\[ C_{\pm} := \{ \lambda \in \mathbb{C} : \lambda = \pm |1 + e^{i\varphi_1} + e^{i\varphi_2}|/3, \varphi_1, \varphi_2 \in [0, 2\pi] \}. \]
The curves \( C_{\pm} \) coincide with the intervals \([0, 1]\) and \([-1, 0]\), respectively, whence \( \text{sp} \Delta_G = [-1, 1] \).

Let now \( v \) be a \( \mathbb{Z}^2 \)-periodic potential and set \( v_j := v(x_j) \) for \( j = 1, 2 \). A calculation similar to example 24 yields that the spectral curves of the Schrödinger operator \( \mathcal{H}_G := \Delta_G + v I \) are
\[ \left\{ \lambda \in \mathbb{C} : \lambda = \frac{1}{2} \pm \frac{\sqrt{(v_1 - v_2)^2 + 4 \mu(\varphi_1, \varphi_2)}}{2(v_1 + v_2)} \right\}, \]
where
\[ \mu(\varphi_1, \varphi_2) := |1 + e^{i\varphi_1} + e^{i\varphi_2}|^2/9, \quad \varphi_1, \varphi_2 \in [0, 2\pi]. \]

Hence, as in example 24, \( \text{sp}_{\text{ess}} \mathcal{H}_G = \text{sp} \mathcal{H}_G \) is given by the union (25).

Let finally \( v \) be a slowly oscillating potential on \( X \). Since the image of the function \( \mu \) is the interval \([0, 1]\), the essential spectrum of the Schrödinger operator on the honeycomb graph \( \Gamma \) is given by formulae (26) and (27). If the condition (28) holds, then a gap \( (a_G, b_G) \) occurs in the essential spectrum of \( \mathcal{H}_G \).

7. A three-particle problem

Let \( \Gamma := (X, \sim) \) be a \( \mathbb{Z}^n \)-periodic discrete graph. We consider the Schrödinger operator
\[ \mathcal{H}u := \Delta_G \otimes I_X + I_X \otimes \Delta_G + (W_1 I_X) \otimes I_X + I_X \otimes (W_2 I_X) + W_{12} I \]
on $l^2(X \times X)$. This operator describes the motion of two particles with coordinates $x^1, x^2 \in X$ with masses 1 on the graph $\Gamma$ around a heavy nuclei located at the point $x_0 \in X$. Therefore, $\mathcal{H}$ is called a 3-particle Schrödinger operator. In (29), $\Delta_\Gamma$ is again the Laplacian on the graph $\Gamma$, $I_X$ is the identity operator on $l^2(X)$, $I = I_X \otimes I_X$ is the identity operator on $l^2(X \times X)$, $W_1$ and $W_2$ are real-valued functions on $X$ defined by

$$W_j(x^j) = w_j(\rho(x^j, x_0)), \quad j = 1, 2,$$

and $W_{12}$ is a real-valued function on $X \times X$ given by

$$W_{12}(x^1, x^2) = w_{12}(\rho(x^1, x^2)).$$

Here $\rho$ denotes the given metric on $X$, and $w_1, w_2$ and $w_{12}$ are functions on $[0, \infty)$ which satisfy

$$\lim_{z \to \infty} w_1(z) = \lim_{z \to \infty} w_2(z) = \lim_{z \to \infty} w_{12}(z) = 0.$$

Clearly, $\mathcal{H}$ is a band operator on $l^2(X \times X)$. We are going to describe its essential spectrum via formula (24), for which we need the limit operators of $\mathcal{H}$ and their spectra. Note that the spectrum of the Laplacian $\Delta_\Gamma$ depends on the structure of the graph $\Gamma$ and that this spectrum has a band structure ($\sigma$ is the union of closed intervals). In examples 24 and 25 we had $\sigma(\Delta_\Gamma) = [−1, 1]$.

We agree upon the following notation. For non-empty subsets $E, F$ of $\mathbb{R}$, we set

$$E + F := \{z \in \mathbb{R} : z = x + y, x \in E, y \in F\}$$

denote their algebraic sum and set $2E := E + E$.

Let $g = (g^1, g^2) : \mathbb{N} \to \mathbb{Z}^n \times \mathbb{Z}^n$ be a sequence tending to infinity. We have to distinguish the following cases (all other possible cases can be reduced to these cases by passing to suitable subsequences of $g$):

**Case 1.** The sequence $g^1$ tends to infinity, whereas $g^2$ is constant. Then the limit operator $\mathcal{H}_g$ of $\mathcal{H}$ is unitarily equivalent to the operator

$$\mathcal{H}_2 := \Delta_\Gamma \otimes I_X + I_X \otimes (\Delta_\Gamma + W_2 I_X).$$

(30)

**Case 2.** Here $g^2$ tends to infinity and $g^1$ is constant. Then the limit operator $\mathcal{H}_g$ of $\mathcal{H}$ is unitarily equivalent to the operator

$$\mathcal{H}_1 := (\Delta_\Gamma + W_1 I_X) \otimes I_X + I_X \otimes \Delta_\Gamma.$$

(31)

**Case 3.** Both $g^1$ and $g^2$ tend to infinity. There are two subcases:

**Case 3a.** The sequence $g^1 - g^2$ tends to infinity. In this case the limit operator is the free discrete Hamiltonian

$$\Delta_\Gamma \otimes I_X + I_X \otimes \Delta_\Gamma$$

the spectrum of which is $2\sigma(\Delta_\Gamma)$.

**Case 3b.** The sequence $g^1 - g^2$ is constant. Then the limit operator $\mathcal{H}_g$ of $\mathcal{H}$ is unitarily equivalent to the operator of interaction

$$\mathcal{H}_{12} := \Delta_\Gamma \otimes I_X + I_X \otimes \Delta_\Gamma + W_{12} I_X.$$ (32)
Note that the operators $\mathcal{H}_1$, $\mathcal{H}_2$ and $\mathcal{H}_{12}$ are invariant with respect to shifts by elements of the form $(0, g)$, $(g, 0)$ and $(g, g)$ of $\mathbb{Z}^n \times \mathbb{Z}^n$, respectively. It follows from proposition 17 that these operators do not possess discrete spectra. From formula (24) we further conclude

$$\text{sp}_{\text{ess}} \mathcal{H} = \text{sp} \mathcal{H}_1 \cup \text{sp} \mathcal{H}_2 \cup \text{sp} \mathcal{H}_{12}. \quad (33)$$

The following proposition is well known. For a proof see [36], theorem VIII.33 and its corollary.

**Proposition 26.** Let $A \in \mathcal{L}(H)$ and $B \in \mathcal{L}(K)$ be bounded self-adjoint operators on Hilbert spaces $H, K$. Then

$$\text{sp}(A \otimes I_K + I_H \otimes B) = \text{sp} A + \text{sp} B.$$  

In our setting, this proposition implies that

$$\text{sp} \mathcal{H}_2 = \text{sp} \Delta_{\Gamma} + \text{sp}(\Delta_{\Gamma} + W_2 I_X).$$

Since the Schrödinger operator $\Delta_{\Gamma} + W_2 I_X$ is a compact perturbation of the Laplacian $\Delta_{\Gamma}$, one has

$$\text{sp}_{\text{ess}}(\Delta_{\Gamma} + W_2 I_X) = \text{sp} \Delta_{\Gamma} \cup \{ \lambda^{(2)}_k \}_{k=1}^\infty,$$

where $\{ \lambda^{(2)}_k \}_{k=1}^\infty$ is the sequence of the eigenvalues of $\Delta_{\Gamma} + W_2 I_X$ which are located outside the spectrum of $\Delta_{\Gamma}$. Thus,

$$\text{sp} \mathcal{H}_2 = 2 \text{sp} \Delta_{\Gamma} + \bigcup_{k=1}^\infty \{ \lambda^{(2)}_k + \text{sp} \Delta_{\Gamma} \}.$$  

In the same way one finds

$$\text{sp} \mathcal{H}_1 = 2 \text{sp} \Delta_{\Gamma} + \bigcup_{k=1}^\infty \{ \lambda^{(1)}_k + \text{sp} \Delta_{\Gamma} \},$$

where the $\lambda^{(1)}_k$ run through the points of the discrete spectrum of $\Delta_{\Gamma} + W_1 I_X$ which are located outside the spectrum of $\Delta_{\Gamma}$.

Recall that in examples 24 and 25, $\text{sp} \Delta_{\Gamma} = [-1, 1]$. Hence, in the context of these examples,

$$\text{sp} \mathcal{H}_j = [-2, 2] \bigcup_{k=1}^\infty [\lambda^{(j)}_k - 1, \lambda^{(j)}_k + 1].$$

One can also give a simple estimate for the location of the spectrum of $\mathcal{H}_{12}$ by means of the following well-known result (see, e.g., [23], p 357).

**Proposition 27.** Let $A$ be a bounded self-adjoint operator on the Hilbert space $H$. Then $[a, b] \subseteq \text{sp} A \subseteq [a, b]$ where

$$a := \inf_{\|h\|=1} \langle Ah, h \rangle, \quad b := \sup_{\|h\|=1} \langle Ah, h \rangle.$$  

This observation implies the following inclusions for the spectra of the operators $\mathcal{H}_1$, $\mathcal{H}_2$ and $\mathcal{H}_{12}$. For $j = 1, 2$ one has

$$2 \text{sp} \Delta_{\Gamma} \subseteq \text{sp} \mathcal{H}_j \subseteq 2 \text{sp} \Delta_{\Gamma} + \left[ \inf_{x \in X} W_j(x), \sup_{x \in X} W_j(x) \right],$$

whereas

$$2 \text{sp} \Delta_{\Gamma} \subseteq \text{sp} \mathcal{H}_{12} \subseteq 2 \text{sp} \Delta_{\Gamma} + \left[ \inf_{y \in X \times X} W_{12}(y), \sup_{y \in X \times X} W_{12}(y) \right].$$
In the context of examples 24 and 25, these inclusions specify to
\[-2, 2 \subseteq \text{sp} \mathcal{H}_j \subseteq \left[ -2 + \inf_{x \in X} W_j(x), 2 + \sup_{x \in X} W_j(x) \right],
\[-2, 2 \subseteq \text{sp} \mathcal{H}_{12} \subseteq \left[ -2 + \inf_{x \in X \times X} W_{12}(x), 2 + \sup_{x \in X \times X} W_{12}(x) \right].

Thus, by theorem 22,
\[\text{sp}_{\text{ess}} \mathcal{H} \subseteq [m - 2, M + 2],\]
where
\[m := \min \left\{ \inf_{x \in X} W_1(x), \inf_{x \in X} W_2(x), \inf_{x \in X \times X} W_{12}(x) \right\},\]
\[M := \max \left\{ \sup_{x \in X} W_1(x), \sup_{x \in X} W_2(x), \sup_{x \in X \times X} W_{12}(x) \right\}.

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