MID-INFRARED SPECTROSCOPY OF LENSED GALAXIES AT 1 < z < 3: THE NATURE OF SOURCES NEAR THE MIPS CONFUSION LIMIT

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ABSTRACT

We present Spitzer IRS mid-infrared spectra for 15 gravitationally lensed, 24 μm-selected galaxies, and combine the results with four additional very faint galaxies with IRS spectra in the literature. The median intrinsic 24 μm flux density of the sample is 130 μJy, enabling a systematic survey of the spectral properties of the very faint 24 μm sources that dominate the number counts of Spitzer cosmological surveys. Six of the 19 galaxy spectra (32%) show the strong mid-IR continua expected of AGNs; X-ray detections confirm the presence of AGNs in three of these cases, and reveal AGNs in two other galaxies. These results suggest that nuclear accretion may contribute more flux to faint 24 μm-selected samples than previously assumed. Almost all the spectra show some aromatic (PAH) emission features; the measured aromatic flux ratios do not show evolution from z = 0. In particular, the high signal-to-noise mid-IR spectrum of SMM J163554.2+661225 agrees remarkably well with low-redshift, lower luminosity templates. We compare the rest-frame 8 μm and total infrared luminosities of star-forming galaxies, and find that the behavior of this ratio with total IR luminosity has evolved modestly from z = 2 to z = 0. Since the high aromatic-to-continuum flux ratios in these galaxies rule out a dominant contribution by AGNs, this finding implies systematic evolution in the structure and/or metallicity of infrared sources with redshift. It also has implications for the estimates of star-forming rates inferred from 24 μm measurements, in the sense that at z ~ 2, a given observed frame 24 μm luminosity corresponds to a lower bolometric luminosity than would be inferred from low-redshift templates of similar luminosity at the corresponding rest wavelength.

Subject headings: galaxies: evolution — galaxies: high-redshift — infrared: galaxies

1. INTRODUCTION

The Spitzer Space Telescope (Werner et al. 2004) has been tremendously successful at detecting star-forming galaxies and active galactic nuclei (AGNs) at z > 1 by their emission in the observed-frame 24 μm band. Down to the confusion limit of ~50 μJy at 24 μm (Dole et al. 2004), the MIPS instrument (Rieke et al. 2004) can detect luminous infrared galaxies (LIRGs)9 out to z ~ 2, ultraluminous infrared galaxies (ULIRGs)10 out to z ~ 3, and hyperluminous infrared galaxies (HLIRGs)11 out to even higher redshifts.

Detectability at z > 0.7 is increased by strong aromatic emission features, at rest-frame 6–12 μm, that pass through the 24 μm band. Indeed, photometric redshifts place ~25%–30% of the faint 24 μm sources at redshifts above 1.4 (Le Floc’h et al. 2005; Pérez-González et al. 2005; Caputi et al. 2006; Wang 2006).

Spectra have been obtained with the Infrared Spectrograph (IRS; Houck et al. 2004) for extremely high-luminosity galaxies, (e.g., Houck et al. 2005; Yan et al. 2007; Menéndez-Delmestre et al. 2007). However, since spectral properties depend strongly on luminosity, results obtained for hyperluminous galaxies may have limited applicability for the bulk of the IR-detected population. Spectra of lower luminosity galaxies have been obtained only in a few cases (e.g., Lutz et al. 2005; Teplitz et al. 2007).

Numerous authors (e.g., Caputi et al. 2006; Choi et al. 2006; Reddy et al. 2006) have used the observed 24 μm band as an estimator of star formation rates. Since most of the infrared power is radiated redward of the observed band, the 24 μm diagnostic must be calibrated using low-redshift templates (e.g., Chary & Elbaz 2001; Dale & Helou 2002; Lagache et al. 2004; Brandl et al. 2006; Armus et al. 2007). The central assumption in these works is that the spectra and spectral energy distributions of high-z galaxies are matched accurately by the low-redshift templates. For 0.7 < z < 3, the 24 μm band fluxes are strongly influenced by aromatic band emission. Evolution over time either in the behavior of these aromatic bands, in metallicity, or in the geometry and radiative transfer within these infrared sources, could all undermine this central assumption.

To probe these possibilities, we have been obtaining IRS spectra for intrinsically faint 24 μm-selected galaxies at 1 ≤ z ≤ 3. Our targets are strongly lensed by the gravitational potential of foreground clusters of galaxies, such that their fluxes are amplified by factors of 3–25. Although these galaxies have observed 24 μm flux densities of ~1 mJy, sufficient for useful, efficient spectroscopy with IRS, their intrinsic luminosities are down in the LIRG to ULIRG range.
In this paper we present IRS spectra for the first 15 galaxies in the sample. The median observed flux density is $f_{\nu}(24 \, \mu m) = 0.82 \, mJy$, but the median intrinsic flux density is nearly an order of magnitude lower. We compare their aromatic feature flux ratios, and the ratios of their X-ray, aromatic, and total infrared luminosities to those of low-redshift LIRG and ULIRG samples. We examine in detail the high-quality spectrum of SMM J163554.2+661225, a LIRG at $z = 2.516$, to measure the sources in four additional objects from the literature (Lutz et al. 2005; Teplitz et al. 2007) to provide a first systematic look at the IRS-measured properties of the faint $24 \, \mu m$ galaxy population.

We assume $\Omega_m = 0.265$, $\Omega_{\Lambda} = 0.735$, and $H_0 = 72 \, km \, s^{-1} \, Mpc^{-1}$ (Spergel et al. 2003; Freedman et al. 2001).

### Table 1: The IRS Lensed Sample

| Source                  | R.A. (J2000) | Decl. (J2000) | Submillimeter Name | Reference | PID | $z$ | $z$ Reference | Amplitude |
|-------------------------|--------------|---------------|-------------------|-----------|-----|-----|---------------|------------|
| MS0451                 | 04 54 07.12  | -03 00 36.4   | SMM J040429+4658  | A         | 2   | 1.95| 1             | 2.5        |
| A851a                   | 09 42 54.55  | +46 58 44.3   | SMM J13115+01208  | B         | 2   | 1.537| 1             | 8.5 ± 2.3  |
| A1689b                  | 13 11 28.04  | -01 19 18.7   | SMM J140105+025223.5 | C       | 12  | 2.565| 3             | 3.5 ± 0.5a |
| A2218a                  | 16 35 54.18  | +66 12 24.8   | SMM J163554.2+661225 | D       | 1   | 2.516| 4             | 22 ± 2     |
| A2218b                  | 16 35 55.16  | +66 11 50.8   | SMM J163554.2+661150 | E       | 1   | 1.034| 5             | 6.1        |
| A2218c                  | 16 35 59.12  | +66 11 47.4   | SMM J163554.2+661150 | F       | 2   | 0.97 | 1             | 6.7        |
| A2219a                  | 16 40 19.50  | +46 44 00.5   | SMM J164034+46440 | G       | 1   | 2.03 | 1             | 3.6        |
| A2261a                  | 17 22 21.00  | +32 07 05.5   | SMM J172235+3207  | F         | 12  | 1 < z < 3 | 1             | 3.3        |
| A2390a                  | 21 53 33.24  | +17 42 10.6   | SMM J182235+3207  | F         | 12  | 2.585| 6             | 10.8       |
| A2390b                  | 21 53 34.43  | +17 42 21.6   | SMM J182235+3207  | F         | 12  | 0.913| 7             | 3.35       |
| A2390c                  | 21 53 34.54  | +17 42 03.0   | “Source 14”       | G         | 2   | 0.91 | 8             | 10.4       |
| AC 114a                 | 22 58 49.79  | -34 48 47.0   |                    | G         | 2   | 1.47 | 1             | 9.7        |
| AC 2667a                | 23 51 40.00  | -26 04 52.0   |                    | G         | 2   | 1.034| 9             | 17         |

### Notes
- Col. (1): Source name.
- Cols. (2)–(3): Coordinates: units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.
- Col. (4): Submillimeter name, if source has a submillimeter detection in the literature.
- Col. (5): References for detected submillimeter counterpart: (A) Smail et al. 2002; (B) Knudsen 2004 and 2008; (C) Smail et al. 2000 and Ivison et al. 2000; (D) Kneib et al. 2004; (E) Knudsen et al. 2006; (F) Chapman et al. 2002; (G) Cowie et al. 2002.
- Col. (6): Redshift and redshift reference: (1) redshift determined from our IRS spectra; (2) redshift based on Keck spectra (J. Richard et al. in prep.); (3) redshift based on Keck spectra (J. Richard et al. in prep.).
- Col. (7): Spitzer program in which IRS spectra were obtained: (1) PID 82; (2) PID 30775; (5) PID 252; (9) PID 3241.
- Cols. (7)–(8): Redshift and redshift reference: (1) redshift determined from our IRS spectra; (2) redshift based on Keck spectra (J. Richard et al. in prep.); (3) from Frayer et al. 1999; (4) Kneib et al. 2004; (5) Ebbels et al. 1998; (6) Keck spectra presented in the Appendix; (7) Pello et al. 1991; (8) Frye & Broadhurst 1998; (9) Sand et al. 2005; (10) Teplitz et al. 2007; (11) Frayer et al. 1998 and Lutz et al. 2005.
- Col. (9): Amplitude estimated from the lensing models. Values from the literature are footnoted; un-footnoted values are new.

The amplification for this source has been the subject of several papers; Downes & Solomon (2003) and Motohara et al. (2005) argue that an intervening cluster galaxy boosts the amplification to $\sim 25$. Smail et al. (2005) and Smith et al. (2005) find that the intervening galaxy is a dwarf with small velocity dispersion and thus contributes only modestly to the total amplification, which they estimate as $3.5 \pm 0.5$. We note the controversy and adopt the Smail et al. (2005) amplification.

2. SAMPLE SELECTION

Our goal was to select sources with intrinsic flux densities close to, or even below, the MIPS $24 \, \mu m$ confusion limit. We first obtained deep $24 \, \mu m$ images of galaxy clusters, generally identified as being massive through high X-ray luminosity. From these $24 \, \mu m$ images, we selected sources within $1.5'$ of the cluster center that met the following criteria:

1. Had observed $24 \, \mu m$ flux densities above 0.4 mJy (to make efficient use of IRS).

2. Had faint, irregular, or arclike HST counterparts. This excluded galaxies at the cluster redshift.12

3. Had known spectroscopic or probable photometric redshifts $\geq 1.0$.

4. Had lensing amplification factors calculated to be greater than 3. This restricted selection to clusters with good lensing models. Since amplification depends on the source redshift, if the redshift was unknown, we calculated the amplification for $1 < z < 3$ and required amplifications factors greater than 3 over that range.

Since submillimeter sources in lensing clusters have been studied extensively, with published amplifications and redshifts, we prioritized sources in the target list with published submillimeter detections. Our sample contains 15 galaxies, targeted in GTO campaigns 82 and 30775. Eight of these galaxies have submillimeter detections in the literature, as referenced in Table 1. Figure 1 shows optical-band images of the galaxies.

We supplement the sample with four sources from the literature. We use three of the four sources from Teplitz et al. (2007) A2219a and A2261a lack HST imagery; we instead used ground-based F606W images from the Steward Observatory 90 inch telescope. Both galaxies are known submillimeter sources (see Table 1), and can thus be added to the sample despite the lack of high-resolution optical images.

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with $f_{\nu}(24\,\mu m) = 0.13-0.2\,mJy$; these are unlensed galaxies at $1 < z < 3$ with very long (9–12 hr) IRS exposures, and are not known submillimeter sources. We also use the lensed $z = 2.81$ submillimeter source in A370 from Lutz et al. (2005)\textsuperscript{13} with $f_{\nu}(24\,\mu m) = 1.36\,mJy$.

Source coordinates, redshifts, amplifications, and flux densities for our sample and for the four literature sources are all listed in Tables 1 and 2.

3. OBSERVATIONS AND DATA REDUCTION

This paper is based on a combination of imaging with MIPS, spectroscopy with IRS, archival imaging from Chandra, and submillimeter photometry from the literature. Exposure times are summarized in Table 3.

3.1. Spitzer IRS Spectroscopy

Low-resolution IRS spectra were obtained as part of Spitzer GTO programs 82 and 30775 (PI G. Rieke). Data reduction used a package developed by D. Fadda, which has already been successfully used to reduce low-resolution IRS spectra of faint high-$z$ sources (Yan et al. 2007). Residual background, rogue pixels\textsuperscript{14} (pixels with dark current values abnormally high and variable with time), and cosmic-ray hits have to be corrected to obtain an optimal reduction of the data. The package considers all the bidimensional frames produced by the IRS/SSC pipeline. First,

\textsuperscript{13} The two other sources from those papers, “2-x” from Teplitz et al. (2007) and the Abell 2125 source from Lutz et al. (2005) have insufficient signal-to-noise for our analysis.

\textsuperscript{14} See http://ssc.spitzer.caltech.edu/irs/roguepixels.
Table 2

| Source        | $f(3.5 \, \mu m)$ (mJy) | $f(4.5 \, \mu m)$ (mJy) | $f(5.7 \, \mu m)$ (mJy) | $f(8.0 \, \mu m)$ (mJy) | $f(24 \, \mu m)$ (mJy) | $f(70 \, \mu m)$ (mJy) | $L(TIR)$ (L$_{\odot}$) | Template |
|---------------|-------------------------|-------------------------|-------------------------|-------------------------|-------------------------|-------------------------|-------------------------|----------|
| MS0451        | 32 ± 2                  | 49 ± 2                  | 52 ± 3                  | 36 ± 3                  | 1.32                    | 6.3E–11                 | 5.7 ± 1.6               | <2.2 × 10$^{13}$ CE*  |
| A851a         | 29 ± 3                  | 44 ± 3                  | 58 ± 3                  | 45 ± 4                  | 0.63                    | 8.4E–11                 | ...                     | 2.7 (2.3–3.6) × 10$^{11}$ DH |
| A1689a        | 36 ± 2                  | 35 ± 2                  | 28 ± 2                  | 66 ± 4                  | 0.58                    | ND                      | ...                     | ND       |
| A1689b        | CR                      | CR                      | CR                      | 0.32                    | 2.4E–10                 | ...                     | 1.6 (0.43–2.4) × 10$^{11}$ DH |
| A1835a        | 88 ± 2                  | 105 ± 2                 | 136 ± 3                 | 124 ± 5                 | 0.99                    | 1.4E–10                 | <13                    | 6.8 (4.2–11) × 10$^{12}$ DH |
| A2218a        | 67 ± 2                  | 70 ± 2                  | 92 ± 6                  | 105 ± 10                | 1.16                    | 2.0E–10                 | <7                     | 7.6 (5.7–9.5) × 10$^{11}$ DH |
| A2218b        | 182 ± 4                 | 147 ± 3                 | 107 ± 3                 | 107 ± 5                 | 1.67                    | ND                      | 7.4 ± 1.5              | 4.8 (2.8–5.3) × 10$^{11}$ DH |
| A2218c        | 99 ± 3                  | 79 ± 3                  | 67 ± 2                  | 523 ± 11                | 0.50                    | ND                      | 10 ± 1                 | 1.3 (0.88–2.0) × 10$^{11}$ Mk 231 |
| A2219a        | CR                      | CR                      | CR                      | 0.82                    | 7.1E–11                 | ...                     | 2.9 (2.7–3.4) × 10$^{12}$ DH |
| A2261a        | 142 ± 4                 | 138 ± 4                 | 117 ± 3                 | 156 ± 6                 | 0.58                    | ND                      | <3.2                   | ...      |
| A2390a        | 39 ± 3                  | 59 ± 3                  | 90 ± 3                  | 196 ± 7                 | 0.83                    | 5.8E–11                 | <7.6                   | ...      |
| A2390b        | 462 ± 8                 | 334 ± 6                 | 258 ± 5                 | 199 ± 7                 | 0.88                    | ND                      | 5.3 ± 1.6              | 1.7 (1.1–2.5) × 10$^{11}$ Mk 231 |
| A2390c        | CR                      | CR                      | CR                      | 0.60                    | ND                      | <8                     | 8.2 (5.5–11) × 10$^{10}$ DH |
| AC114a        | 91 ± 3                  | 97 ± 3                  | 78 ± 3                  | 76 ± 4                  | 0.41                    | 1.4E–10                 | ...                     | ...      |
| A2667a        | 466 ± 6                 | 349 ± 5                 | 269 ± 4                 | 320 ± 8                 | 1.52                    | ND                      | ...                     | ...      |

Notes.—"ND" indicates a nondetection; ellipses indicate no data. Col. (1): Source name. Cols. (2)–(5): IRAC photometry, in mJy. Formal errors from the photometry are quoted, although field crowding probably and aperture corrections probably introduce additional 5%–10% errors. The photometry for A2218a is superior, since it is multiply imaged (with different crowding). "CR" indicates that crowding was too severe to quote photometry. Col. (6): Observed 24 μm flux density, in mJy, from DAOPHOT PSF fitting. Error bars from DAOPHOT are overly optimistic and not quoted; we use error bars of 0.1 mJy for template fitting. Col. (7): Aromatic–to–mid-IR flux ratio, as defined in § 4.2, for the z > 1.4 sample. Col. (8): Observed 70 μm flux density, in mJy, from aperture photometry. Error bars are dominated by the sky. Col. (9): Total infrared luminosity, our best–fit and the range of acceptable templates, in L$_{\odot}$. Col. (10): Best-fitting template: Chary & Elbaz (2001) or Dale & Helou (2002).

A Mk 231 template may also fit.

b Using amplification factor of 1.3.

c L(TIR) cannot be usefully constrained because the redshift is unknown.

Table 3

| Cluster       | t(24 μm) (ks) | t(3.6 μm) (ks) | t(IRS, LL) (ks) | t(Chandra) (ks) |
|---------------|--------------|---------------|----------------|-----------------|
| MS 0451       | 3.68         | 2.40          | 3.66$^a$       | 53              |
| A851          | 2.78         | 2.40          | 6.34$^a$       | 0               |
| A1689         | 2.78         | 2.40          | 7.31$^a$       | 40              |
| A1835         | 2.78         | 6.00          | 3.66$^a$; 3.66$^a$ | 29             |
| A2218         | 2.78         | 2.40          | 3.66$^a$       | 58              |
| A2261         | 2.77         | 2.40          | 5.89$^a$; 7.31$^a$ | 33             |
| A2390         | 2.77         | 2.40          | 3.66$^a$; 5.61$^a$ | 109            |
| AC 114        | 2.77         | 2.40          | 7.31$^a$       | 70              |
| A2667         | 2.77         | 2.40          | 1.95$^a$       | 9.2             |

Notes.—Exposure times, in kiloseconds, for Spitzer MIPS 24 μm, IRAC 3.6 μm, and IRS Long-Low grating. The 24 μm exposure time is the median (per-pixel) exposure time within the central 4’’×4’’ box.

$^a$ Both LL1 and LL2 were given this exposure time.

$^b$ Only order LL1.

background and noise images are produced by co-adding frames after masking target and serendipitous spectra on each frame. (The background is better estimated for the PID 30775 observations where both LL1 and LL2 orders provide redundant background measurements, compared to the PID 82 observations using only LL1, where the background can only be determined by different nodding and thus potentially suffers from contamination by faint sources.)

A robust statistical estimator (biweight) is used to minimize the effects of deviant pixels on the co-added value. This part is iterative in order to allow manual identification of sources. Rogue pixels are then identified by computing the dispersion of the noise around every pixel and flagging pixels which are >5 σ deviant from the mean local value. Spectral extraction is done optimally (e.g., by using the PSF profile to weight the spectrum), taking into account the spectral distortion, and also rejecting pixels affected by cosmic rays. Finally, the spectra obtained at each position are co-added.

We downloaded and reduced one archival IRS spectrum, from program 3453 (PI P. van der Werf). This is SMM J163555.2+661238, image "A" in Kneib et al. (2004), one of three images of the same triply-imaged submillimeter galaxy in Abell 2218. We obtained an IRS spectrum of the brightest image of this galaxy, SMM J163554.2+661225 (image "B" in Kneib et al. 2004); its 24 μm flux is 1.8 times brighter than the SMM J163555.2+661238. In the archival spectrum, the source was miscentered across the short axis of the slit, which reduced throughput. Because its signal-to-noise ratio is low, we use the archival spectrum only to confirm
Fig. 2.—IRS spectra. Bottom x-axes show observed wavelength in $\mu$m; top x-axes show rest wavelength ($\mu$m) if redshift is known; y-axes plot observed flux density (mJy). Spectra obtained for the same object in different programs are plotted separately, with program 82 plotted first. Vertical dotted lines show the positions of expected aromatic and fine structure emission features. Note that artifacts are present at $\lambda > 34\mu$m.
the detection of features, and use the spectrum from our program for all analysis.

The IRS spectra are presented in Figure 2. Spectra obtained in programs 82 and 30775 for the same object are plotted separately; observations from the latter program have superior background subtraction. Flux densities are as observed, without any correction for gravitational amplification. When redshifts were known from the literature, or are evident from aromatic emission features in the IRS spectrum, then the upper $x$-axis shows rest wavelength. Only one source (A2261a) lacks a redshift. Redshifts are listed in Table 1, with a reference if the redshift is from the literature.

The rest-frame spectra were run through PAHFIT (Smith et al. 2007) to fit the continuum and aromatic (PAH) features simultaneously. Measured aromatic fluxes are reported in Table 4. We allowed PAHFIT to fit the silicate absorption (e.g., we did not specify “NO_EXTINCTION”). Fixing the silicate optical depth at zero changes the best-fit 7.7 $\mu$m feature fluxes by 4% (median), the best-fit 8.6 $\mu$m feature fluxes by 8% (median), and the best-fit 11.3 $\mu$m feature fluxes by 11% (median).

We report fluxes in Table 4 for features in which the measured flux is at least twice the uncertainty calculated for that line. We add back in a few features, mostly at 11.3 $\mu$m, which are obviously detected even though the error bars are large.

3.2. Spitzer MIPS 24 and 70 $\mu$m Photometry

MIPS photometry-mode images at 24 and 70 $\mu$m were obtained through Spitzer GTO program 83 (PI G. Rieke). Photometry is given in Table 2, and exposure times are given in Table 3. The data were reduced and mosaicked using the Data Analysis Tool (Gordon et al. 2006) with a few additional processing steps (Egami et al. 2006b).

Photometry at 24 $\mu$m was obtained by PSF fitting, using the IRAF implementation of the DAOPHOT task allstar. The PSF was created empirically using all available images from program
| Source       | 6.2 | σ  | 7.7 | σ  | 8.3 | σ  | 8.6 | σ  | 11.3 | σ  | 12.0 | σ  | 12.6 | σ  | 13.6 | σ  | 14.2 | σ  | 16.4 | σ  | 17  | σ  |
|--------------|-----|----|-----|----|-----|----|-----|----|------|----|------|----|------|----|------|----|------|----|------|----|-----|----|
| MS 0451a..... |    |    | 28  | 6  | 4.8 | 2  | 4.3 | 2  | 5.0  | 2  | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... |
| A851a......... | 2.8 | 0.6| 15.6| 0.8| 2.4 | 1  | 2.7 | 0.8| ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... |
| A1689a........ |    |    | 12  | 5  |      |    |      |    | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... |
| A1689b........ | 6.5 | 1.1| 21  | 3  | 3.8 | 1  | 4.0 | 2  | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... |
| A1835a........ | 13  | 2  | 39  | 10 | 7.3 | 2  | 7.8 | 2  | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... |
| A2218a........ | 18.3| 0.9| 66  | 3  | 8.2 | 1  | 13.5| 1  | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... | ...  | ... |
| A2218b........ |    |    | 65  | 3  | 4.8 | 1.7| 16.3| 2  | 20   | 3  | 3.8  | 1.2| 10.8| 1.3| 1.7  | 0.8| 14   | 5  |
| A2218c........ |    |    | 14.9| 0.9|     |    | 1.5 | 0.8| 1.3  | 0.4|      |    |      |    |      |    |      |    |      |    |      |    |
| A2219a........ |    |    | 19.3| 1.3|     |    | 4.6 | 1  | 5.6  | 2  |      | ...| ...  | ...| ...  | ...| ...  | ...| ...  | ...| ...  | ...|
| A2390a........ | 1.1 | 1  | 12.4| 3  |     |    | 5.8 | 2  | ...  | ...| ...  | ...| ...  | ...| ...  | ...| ...  | ...| ...  | ...| ...  | ...|
| A2390b........ |    |    |     |    |     |    |     |   | 14   | 10| 6.1  | 2.3|      |   |      |   |      |   |      |   |      |    |
| A2390c........ |    |    | 35  | 11 |     |    | 11  | 8  | 17   | 12| 6.3  | 1.5|      |   |      |   |      |   |      |   |      |    |
| AC 114a........ | 6.5 | 0.5| 23  | 2  |     |    | 7.1 | 1  | 9.2  | 1  | 4.2  | 2  | ...  | ...| ...  | ...| ...  | ...| ...  | ...| ...  | ...|
| A2667a......... |    |    | 102 | 11 | 12  | 3  | 22  | 4  | 29   | 7  | 9.1  | 3  | 16   | 3 | 5.6  | 1.9|      |   |      |   |      |    |

Notes.— Aromatic dust emission feature fluxes and uncertainties, in units of $10^{-15}$ erg s$^{-1}$ cm$^{-2}$, measured with PAHFIT. Blank cells are nondetections. Column headers give the central wavelength, in μm, of each feature. For features comprising multiple components (such as 7.7 μm), total flux and uncertainty were calculated using PAHFIT’s “main feature” function. Ellipses indicate no coverage; blank indicates nondetections.
3.3. Spitzer IRAC Photometry

IRAC images at 3.6, 4.5, 5.8, and 8 μm were obtained as part of GTO program 83 (PI G. Rieke) for five of the six clusters; IRAC data for Abell 1835 were obtained through GTO program 64 (PI G. Fazio). Exposure times are listed in Table 3. IRAC images were mosaicked as described by Huang et al. (2004).

Photometry was obtained by aperture photometry with sky subtraction. The high density of cluster galaxies in the cluster cores necessitated use of irregularly shaped polygons for both target and sky apertures. The flux density within these polygonal regions was determined using CIAO\textsuperscript{17} versions 3.2.1 and 3.4.1.1. IRAC photometry is listed in Table 2; the values do not include aperture corrections.

3.4. HST Photometry from WFPC2 and ACS

From the HST archive, we downloaded all publicly available ACS images for the clusters. We used multidrizzle (Koekemoer et al. 2002) to distortion-correct each flat-fielded pipeline image, and then cross-correlated the images to measure the small (<1") coordinate registration offsets with high precision. The measured offsets were then used by multidrizzle for the final mosaicking. The data were photometered using the same technique as for the IRAC bands.

We also used reduced WFPC2 images kindly made available by D. Sand, which are described in Sand et al. (2005). ACS and WFPC2 cutouts of the galaxies in our sample are shown in Figure 1.

3.5. Chandra ACIS Imagery

From the Chandra archive, we downloaded all publicly available ACIS observations of clusters with IRS targets in this sample: observation ID numbers 529 (obs/00529) and 902 (obs/00902) for MS 0451; 1663 (obs/01663), 5004 (obs/05004), and 540 (obs/0540) for A1689; 495 (obs/00495) and 496 (obs/00496) for A1835; 1454 (obs/01454), 1666 (obs/01666), and 553 (obs/00553) for A2218; 896 (obs/00896) for A2219; 5007 (obs/00507) and 550 (obs/00550) for A2261; 4193 (obs/04193), 500 (obs/00500), and 501 (obs/00501) for A2390; and 1562 (obs/01562) for AC114.

We used CIAO (3.3) to update the CTE correction when necessary and remove data obtained during periods of high solar flares. We mosaicked multiple overlapping observations following the CIAO thread "Reprojecting Images: Making an Exposure-corrected Mosaic."

For X-ray detections, fluxes were determined from the individual (not mosaicked) observations as described in Marcillac et al. (2007). X-ray flux upper limits (3 σ) were determined from the mosaicked images and mosaicked exposure maps as in Donley et al. (2005) except that the 90% encircled energy radius was used, and the Monte Carlo sky apertures were taken close to the source (within 30″) in order to sample a representative cluster background.

Chandra X-ray fluxes and upper limits for the 2–8 and 0.5–8 keV bands are reported in Table 5.

3.6. Long-Wavelength Photometry from the Literature

Eight galaxies in our sample, and one galaxy from the extended (literature) sample, have submillimeter detections in the literature; these are listed and referenced in Table 1. Four sources were detected at 15 μm by Infrared Space Observatory (ISO): A851a and A2218b (Barvains et al. 1999), A2390b (Lemonon et al. 1998), and A2390c (Metcalfe et al. 2003). We use this photometry to help determine the total infrared luminosities of the sources.

4. DISCUSSION

4.1. Aromatic Feature Flux Ratios in High-Redshift Galaxies

We now test whether the mid-IR spectra of star-forming galaxies, with their prominent aromatic emission features, show evidence for evolution from z = 0 to high redshift.

Given its high quality, we first examine the IRS spectrum for M82 (Barvains et al. 2005) except that the 90% encircled energy radius was

| Source       | f(2–8 keV)       | f(0.5–8 keV)       |
|--------------|------------------|------------------|
| MS 0451a     | <1.55 x 10^{-15} | <1.30 x 10^{-15} |
| A851a        | …                | …                |
| A1689a       | <7.9 x 10^{-15}  | <4.5 x 10^{-15}  |
| A1689b       | <1.1 x 10^{-14}  | <7.2 x 10^{-15}  |
| A1835a       | <5.8 x 10^{-15}  | <4.4 x 10^{-15}  |
| A2218a       | <5.2 x 10^{-15}  | <2.6 x 10^{-15}  |
| A2218b       | <2.8 x 10^{-15}  | <2.1 x 10^{-15}  |
| A2218c       | <2.4 x 10^{-15}  | <1.6 x 10^{-15}  |
| A2219a       | <3.4 x 10^{-15}  | <3.0 x 10^{-15}  |
| A2261a       | 3.7 x 10^{-15}   | 1.8 x 10^{-15}   |
| A2390a       | 5.1 x 10^{-14}   | 7.6 x 10^{-14}   |
| A2390b       | 4.4 x 10^{-14}   | 2.5 x 10^{-14}   |
| A2390c       | <7.5 x 10^{-15}  | <5.2 x 10^{-15}  |
| AC 114a      | <2.8 x 10^{-15}  | <2.9 x 10^{-15}  |
| A2667a       | <3.2 x 10^{-14}  | 4.1 x 10^{-14}   |

Notes.—X-ray limits are the 3 σ sky background plus any additional counts at the source position, as described in § 3.5.
for the Brandl template, and those reported for the SINGS sample. The strength of the 8.6 μm feature in SMM J163554.2+661225 relative to those at 6.2 and 7.7 μm, is within the 10%–90% variation within the SINGS sample. In addition, the flux ratios of SMM J163554.2+661225 are a close match to those of the Brandl template, and those reported for the SINGS sample (Smith et al. 2007); (2): flux ratios of the average starburst template of Brandl et al. (2006).

Two of the four galaxies from the literature, and three of the galaxies in our program are detected in X-rays: two have no apparent aromatic features (sources A2261a and A2390a); and the other, source A2390b, is interesting in that it contains a luminous X-ray—emitting AGN, yet its spectrum still shows aromatic features. That source emits roughly equal power in the 7.7 μm aromatic feature (scaled from the 11 μm feature) and at 10–30 keV, reminiscent of the more luminous z = 1.15 source CXO GWS J141741.9+522823 discussed by Le Floc’h et al. (2007).

We now consider a second AGN diagnostic, the relative contribution of aromatic versus continuum emission to the mid-IR flux. At low redshift, low aromatic feature equivalent widths have been demonstrated as an effective AGN diagnostic (Armus et al. 2007; Brandl et al. 2006). However, it is extremely difficult to measure equivalent widths accurately for the high-redshift IRS sources, due to the limited wavelength baseline, the fact that the aromatic features have broad wings, and most importantly, imperfect sky subtraction which adds a (positive or negative) pedestal to each spectrum. For example, for SMM J163554.2+661225 (our highest quality spectrum), we cannot measure an accurate equivalent width due to the difficulty in determining the true continuum level. We are able to set a lower limit of >0.4 μm (rest-frame, conservatively assuming a high continuum level), which is a typical value for low-redshift star-forming galaxies. For the other spectra in the sample, equivalent widths are even harder to measure.

Therefore, we instead create an aromatic-to-mid-IR flux ratio, which we define as the ratio of the flux in the 7.7 μm PAH feature (in erg s⁻¹ cm⁻², as fit by PAHFIT) to the MIPS photometry at 24 μm (in Jy, corrected for bandwidth compression to z = 0.). This metric does not suffer as strongly from the difficulty in determining the continuum level. In the low-redshift, star-forming comparison sample we construct in § 4.4, this ratio ranges from (7–32) × 10⁻¹¹. We thus take 7 × 10⁻¹¹ as a dividing line between spectra dominated by star formation (above the value), and spectra with substantial AGN contribution (below the value). This discriminant selects the following mid-IR spectra as having substantial AGN contribution: A370a, MS 0451a, A1689a, A2218c, A2261a, and A2390a. Thus, of 19 galaxies with adequate spectra, six show strong indications of AGN contribution to their mid-IR outputs.

### Table 6

| Lines | Ratio | SINGS | Starburst |
|-------|-------|-------|-----------|
| 7.7 μm/6.2 μm | 3.6 ± 0.2 | 3.6 (1.3–4.8) | 3.9 |
| 7.7 μm/8.6 μm | 4.9 ± 0.4 | 5.7 (4.7–9.0) | 5.1 |
| 6.2 μm/8.6 μm | 1.36 ± 0.1 | 1.5 (1.2–3.0) | 1.3 |

**Notes.**—Col. (1): Line wavelengths (μm); Col. (2): measured flux ratios; Col. (3): median flux ratios and 10%–90% range of variation from the SINGS sample (Smith et al. 2007); Col. (4): flux ratios of the average starburst template of Brandl et al. (2006).

The first AGN diagnostic we examine is X-ray luminosity, plotted in Figure 5 against aromatic luminosity. Most of the sample are X-ray nondetections, at limiting luminosities that rule out X-ray-loud QSOs or bright Seyfert galaxies, although an X-ray-weak or highly obscured AGN could still be present (see Donley et al. 2005; Alonso-Herrero et al. 2006a).

### 4.2. Evidence for Compact Source Accretion

Active galactic nuclei are detectable in multiple ways—by the dust they heat that radiates in the mid-IR; by X-rays emitted from the accretion disk corona; and by high-excitation or broad emission lines. We now examine the evidence for AGN activity in our extended sample.

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18 For A2390b, we scale from the flux in the 11 μm aromatic feature, using the average 7.7/11 μm flux ratio of 3.6 from Smith et al. (2007).
A third indication of AGN activity is a mid-IR spectrum that rises steeply with increasing wavelength. Four spectra (A2261a, A2218c, A1689a, and A370a) show this behavior; none has detected aromatic emission.

A fourth indication of AGN activity is the presence of broad or highly excited lines in an optical spectrum. Of the spectroscopy that has been published, or is in preparation, for our sample (see references in Table 1), three show evidence for AGNs: A1689a has a Keck spectrum showing highly ionized neon (J. Richard et al., in preparation); A2390a has a Keck spectrum (see Appendix) which shows Ly$\alpha$ with FWHM $\sim$ 945 km s$^{-1}$, typical of an AGN narrow-line region; and the spectrum of A370a is reported to contain AGN lines (Lutz et al. 2005).

A fifth indication of AGN activity is a high ratio of [Ne iii] 15.5 $\mu$m to [Ne ii] 12.8 $\mu$m. Solar-metallicity starbursts are not observed to excite this line ratio above unity (Thornley et al. 2000; Rigby & Rieke 2004), whereas the harder continuum of AGNs do. Eight spectra cover rest-frame 12.8 $\mu$m; [Ne ii] is detected in six of these. Six spectra cover rest-frame 15.5 $\mu$m; [Ne iii] is detected in four. Line fluxes are added in Table 7. Of the four galaxies where both lines are detected (A1689a, Ab2218b, A2390b, and A2667a), only in A1689a does the [Ne iii]/[Ne ii] ratio exceed unity: 2.15 $\pm$ 0.76. This high line ratio offers additional evidence that A1689a hosts a luminous AGN. The other, lower line ratios are more difficult to interpret.

The results for these AGN diagnostics are compiled in Table 8. Combining these diagnostics, 8 galaxies out of the sample of 19 have at least one indication of AGN activity (X-ray detection, low aromatic contribution to the mid-IR flux, rising mid-IR spectrum, or optical AGN lines). Five galaxies have two different diagnostics that indicate AGNs. Thus, we find that 42% of the faint 24 $\mu$m sources in this sample have some nuclear activity; the 95% confidence interval is 23%–63%. Further, in 32% of the sample, the AGN strongly affects or dominates the mid-IR spectrum; the 95% confidence interval is 15%–53%. (Confidence intervals were calculated using eq. [26] of Gehrels 1986.)

The high proportion of AGNs in our sample of high-redshift LIRGs and ULIRGs mirrors the general behavior of local IR-luminous galaxies, where AGNs become substantially more common with increasing luminosity (e.g., Lutz et al. 1998). Although the error bars are large given the small sample size, the high proportion (32%) of AGNs sufficiently powerful to dominate the mid-IR spectra may be inconsistent with the frequent assumption that $\sim$100 $\mu$Jy, $z > 1$ sources are dominated by star formation, with AGNs a modest contaminant. Our result, if confirmed with better statistics, has implications for 24 $\mu$m–derived star formation rates in deep surveys, where it is often assumed that the mid-IR outputs of high-redshift, faint galaxies are dominated by strong aromatic features from star formation. It also implies that a higher fraction of the integrated

![Figure 4](https://example.com/figure4.png)

**Fig. 4.** Aromatic flux ratios. The aromatic flux ratios are plotted against observed flux in the 7.7 $\mu$m feature. Fluxes were fit using PAHFIT. Amplifications have not been divided out. Also plotted are the flux ratios of the low-redshift Brandl et al. (2006) starburst template (solid horizontal lines), with the ±1 standard deviations from that sample (dotted horizontal lines).
24 \mu m flux on the sky may arise from accretion than previously thought.

4.3. Total Infrared Luminosity

We now determine the total infrared (TIR) luminosities of the lensed sources, so that we may examine the dependence of aromatic feature emission on total infrared luminosity, and how this behavior may evolve with redshift.

4.3.1. SMM J163554.2+661225

The submillimeter galaxy behind Abell 2218 has exceptionally good photometric coverage from 0.4 to 850 \mu m. Since it is triply-lensed, we increase the signal-to-noise ratio of its multi-band photometry by using fluxes from all three images, weighed by the amplifications given in Kneib et al. (2004). (Bright neighboring galaxies forced us to use only image A and B for the optical and IRAC photometry.) Figure 6 plots this photometry.

We fit redshifted Chary & Elbaz (2001) and Dale & Helou (2002) templates to the 24, 70, 450, and 850 \mu m photometry, corrected for lensing amplification. For each template family, we determine the best-fit template and the range of acceptable templates, then determine \( L(TIR) \) for those templates in a method consistent with the definition given in Table 1 of Sanders & Mirabel (1996) using 12, 25, 60, and 100 \mu m photometry. We estimate the intrinsic total infrared luminosity of SMM J163554.2+661225 as follows:

1. \( 7.6 \pm 1.9 \times 10^{11} L_\odot \) for Dale & Helou (2002) templates.
2. \( 1.3^{+1.0}_{-0.2} \times 10^{12} L_\odot \) for Chary & Elbaz (2001) templates.

As Figure 6 shows, the best-fit Chary & Elbaz (2001) template overpredicts all >24 \mu m points, especially at 70 \mu m. The best-fit Dale & Helou (2002) template is a better match to the shape of the overall infrared spectral energy distribution (SED) and also has high aromatic feature equivalent widths, consistent with the spectrum of the source. Thus, this galaxy’s intrinsic luminosity is that of a LIRG.
4.3.2. Other Galaxies

We adopt a similar approach for the other galaxies with sufficient far-infrared and submillimeter data to provide reasonable constraints. We tried both Dale & Helou (2002) and Chary & Elbaz (2001) templates, fitted to fall just below any upper limits so long as they were compatible with the other constraints, i.e., to give a maximal luminosity. Two galaxies are fit poorly by a star-forming template, and much better by an AGN template, Mrk 231; MS 0451a may be, as well. The ranges of acceptable luminosities and the preferred templates are listed in Table 3. Dale & Helou (2002) templates provided generally better fits to our sample than Chary & Elbaz (2001) templates, as expected given the results of Marcillac et al. (2006) for 0 \( \leq z < 1 \) galaxies, fitting ISO, Spitzer, and VLA photometry.

4.4. 8 μm (Rest)-to-Total Luminosity Ratio

We now consider the aromatic and total infrared (TIR) luminosities of the star-forming galaxies in the lensed sample, and how this ratio behaves as a function \( L(\text{TIR}) \). We compare to higher luminosity star-forming galaxies at \( z \sim 2 \), as well as comparable-luminosity star-forming galaxies at \( z \sim 0.8 \) and \( z \sim 0 \). We interpret the properties of the high-redshift galaxies strictly in terms of properties that are well determined for local analogs, so we can test whether inconsistencies emerge.

4.4.1. The \( z \sim 0 \) Comparison Sample

To probe rest-frame 8 μm and total infrared (TIR) luminosities for local galaxies, we use the SINGS sample (photometry from Dale et al. 2005) supplemented in the ULIRG range by IRAS 09111-1007, 10565+2448, 12112+0305, 14348-1447, 17208-0014, 22491-1808, and Arp 220 (Farrah et al. 2003; Armus et al. 2007), all selected because they appear to be dominated by young stars rather than AGNs.\(^{19} \) We also include IRAS 00262+4251, 01388-4618, 02364-4751, 16474+3430, 23128-5919, and 23365+3604 from Rigopoulou et al. (1999). The strength of the aromatic features relative to the continuum indicates that all these galaxies are also dominated by star formation. For the galaxies from Farrah et al. (2003) and Armus et al. (2007) we determined 8 μm photometry from archival IRAC images. In our reductions, we included a correction for the extended source calibration. For those galaxies in Rigopoulou et al. (1999) we converted the tabulated continuum and aromatic 7.7 μm peak fluxes into 8 μm photometric values by comparing with templates for galaxies in common between Armus et al. (2007): IRAS 12112+0305, 14348-1447, 1525+3609, and 22491-1808, plus Arp 220. The peak-to-peak scatter was a factor of 2 (in agreement with the finding of Chary & Elbaz 2001). We add a variety of LIRGs (Alonso-Herrero et al. 2006b), again avoiding any galaxies with indications of significant levels of AGN power. We took Infrared Astronomical Satellite (IRAS) measurements from Sanders et al. (2003) and computed \( L(\text{TIR}) \) to be consistent with the approach taken in that paper. We define \( L(8 \mu m) \) as being proportional to \( v_j f_j \), and use similar definitions for all luminosities at specific wavelengths. Figure 7 shows the relation between rest-frame \( L(8 \mu m) \) and \( L(\text{TIR}) \) for these \( z \sim 0 \) galaxies.

4.4.2. Results from the Literature at \( z \sim 0.85 \)

We add a sample of infrared galaxies at \( z \sim 0.85 \) from Marcillac et al. (2006). This paper utilizes 15 μm data from ISO, as well

\(^{19} \) IRAS 12112+0305 and 14348-447 are double nuclei sources.
Fig. 6—Continued
as 24 μm data from Spitzer, and, for a subset, radio data from the VLA; the authors found that the galaxies had rest-frame \( L(8\,\mu m)/L(12\,\mu m) \) ratios at the high end of the local galaxy dispersion. We make the assumption that \( L(12\,\mu m) \) correlates with \( L(\text{TIR}) \) in that sample, as it is observed to do well at \( z = 0 \). We computed rest 8 μm luminosities for the members of the sample with \( 0.55 < z < 1.2 \) and adequate 15 μm photometry, and corrected them to the IRAC 8 μm band by comparison with the IRS spectrum of IRAS 2249, a local star-formation-dominated ULIRG. We took \( L(\text{TIR}) \) from the Dale & Helou (2002) template fits (Marcillac et al. 2006). The results are shown in green in Figure 8.

There is no mid-IR spectroscopy for these galaxies, but if they had strong mid-IR contributions from AGNs, we would expect substantial dispersion in the \( L(8\,\mu m) \) to \( L(\text{TIR}) \) ratio. Such a large dispersion is not seen.

4.4.3. Results from the Literature at \( z \sim 2 \)

We now compare the behavior of the 8 μm luminosity versus \( L(\text{TIR}) \) at \( z \sim 2 \) with that seen in local luminous galaxies. We first discuss three results from the literature, and then show how our new measurements complement them.

1. First, we took the stacked \( z \approx 1.9 \) galaxy template of Daddi et al. (2005) and compared the ratio of observed 850 to 24 μm flux densities, referred to the appropriate rest wavelengths for each of the local star-forming galaxies where there is sufficient data to construct an accurate SED template (IRAS 1211, 1434, 17208, 2249, and Arp 220). Excluding Arp 220, for which the ratio is 54, the average value for the other four galaxies is 14, whereas the Daddi et al. (2005) template indicates a value of 6.5. Thus, the template indicates stronger emission at rest 8 μm relative to the submillimeter than is typical of local star-forming ULIRGs.

We also compare with total infrared luminosity derived from submillimeter measurements. Since the rest wavelength for the observed 850 μm point in the Daddi et al. template is 293 μm, we construct the average ratio of \( L(\text{TIR})/L(293\,\mu m) \) for local LIRGs and ULIRGs with \( 4 \times 10^{11} < L(\text{TIR}) < 2.5 \times 10^{12} L_\odot \) and high-quality measurements at 350–450 μm (IRAS 09111, 1211, 1434, and 17208, Arp 220, and NGC 1614). We correct the flux densities for the differing effective filter bandwidths using the IRAS 2259 template. We correct to 293 μm assuming a spectral slope of \(-3.3\), determined from a larger sample of local ULIRG templates. We find

\[
L(\text{TIR}) = (98 \pm 40) L(293\,\mu m)
\]

and then plot the Daddi et al. point as a red square in Figure 7. We did not fit the observations of lower luminosity galaxies, because the scatter at lower luminosities is very large, probably due to varying amounts of cold and warm dust in the infrared-emitting objects.

2. Second, we make use of the stacking analysis carried out by Papovich et al. (2007). They report observed-frame 24 and 70 μm stacked photometry for galaxies that are both undetected in X-rays and have non-power-law IRAC SEDs, and are therefore assumed to be powered predominantly by star formation. We computed the rest 8 and 24 μm flux densities for this category (since they quote values for \( z = 2 \), the rest wavelengths are represented directly by the measurements). We corrected the 8 μm flux density (derived from that observed at 24 μm) to the equivalent for the IRAC band, again using the SED of IRAS 2249. Finally, we compute \( L(\text{TIR}) \) from the relation derived from our local galaxy sample:

\[
L(\text{TIR}) = 86.5 \times L(24\,\mu m, \text{ rest})^\alpha,
\]

where \( \alpha = 0.9758 \).

The rms scatter around this relation is only 0.14 dex. The resulting values are shown in Figure 7 as red diamonds (we omitted the lowest flux density bin from Papovich et al. because the signal-to-noise ratio at 70 μm is too low).
3. Third, we use the results of Yan et al. (2007). We select those galaxies that have $z > 1.5$, and rest-frame $6.2 \mu m$ equivalent widths (measured by Sajina et al. 2007) greater than $0.3 \mu m$. From Armus et al. (2007) this equivalent width threshold should select galaxies dominated by star formation. The galaxies selected (MIPS 289, 8493, 15928, 16144, 22530) are all categorized as “type 1” by Yan et al. (2007), meaning that their spectra contain strong aromatic features. The median rest-frame EW(7.7 $\mu m$) is 2.5 $\mu m$, corroborating that star formation dominates the mid-IR spectra. We converted the tabulated $L(5.8 \mu m)$ for these galaxies into $L(8 \mu m)$ from an average template for our local ULIRG sample (IRAS 2249).

We estimate $L(TIR)$ for these sources using the radio-infrared relation. We describe our adjustment of this relation for use at high redshift in detail in a forthcoming publication (Seymour et al. 2007) but give a short summary here. Yun et al. (2001) have quantified the relation for a large number of local galaxies, finding that the average ratio of flux density at 60 $\mu m$ to 1.4 GHz is 110 independent of luminosity. The results of Yun et al. (2001) have not been $K$-corrected; since the infrared SEDs drop rapidly toward the blue, while the radio spectrum drops slowly, these corrections are a fairly steep function of redshift. They enter at a significant level only for the high-luminosity galaxies, because they tend to be at higher redshift than the lower luminosity ones. Including them increases the indicated ratio to $\sim 170$ at a $L(TIR) = 10^{11} L_\odot$. A second issue is that the radio outputs of LIRGs and ULIRGs often show signs of free-free absorption, which can decrease the outputs at 1.4 GHz. To estimate this effect, we took the 19 star-forming galaxies with adequate data from Condon et al. (1991) and computed an average slope for them of $-0.58$ between 1.49 and 8.44 GHz. Taking the 8.44 GHz value to represent the unabsorbed radio output and assuming the intrinsic slope is $-0.7$ (Condon et al. 1991), we revise the IR/radio ratio downward to 140 to take account of absorption; this value is calculated by extrapolating from 8.44 to 1.4 GHz and comparing results for slopes of $-0.58$ and $-0.7$. The second concern is that the radio-infrared ratio may be a function of redshift (Kovács et al. 2006; Vlahakis et al. 2007). We tested this suggestion by determining a $K$-correction between 1.4 GHz and 850 $\mu m$ for an average of local LIRG and ULIRG templates (A. Alonso-Herrero et al., in preparation) and applying it to a large sample of high-redshift sources. We find that the $K$-correction accurately reproduces the observed behavior with no offset, indicating there is no significant evolution of the relation. We have therefore used a value of 140 for the intrinsic flux density ratio at rest wavelengths to convert the radio measurements to TIR ones. We extrapolated the radio fluxes observed to a rest frequency of 1.4 GHz with a slope of $-0.7$. The redshift usefully makes the observed rest frequency significantly higher than 1.4 GHz, reducing the effects of free-free absorption for the high-$z$ galaxies. To convert to $L(TIR)$, we use the relationship from our local sample, $L(TIR) = 740 \times L(60) ^{\alpha}$, where $\alpha = 0.9427$ with an rms scatter of 0.08 dex.

It is important to point out that both of the corrections we have derived have the effect of reducing the change we find in $L(8 \mu m)/L(TIR)$. That is, using the relation derived by Yun et al. (2001) and applying the proposed change at high redshift (Kovács et al. 2006; Vlahakis et al. 2007) would have both resulted in a much larger change than reported. Therefore, we are taking a conservative approach.

We take radio flux densities for the galaxies observed by Yan et al. (2007) from Condon et al. (2003) as tabulated on the SSC World Wide Web site. Two of the remaining galaxies are detected at low significance as indicated by inspection of the grayscale images for them (MIPS 429 and 506), so we enter rough values of 0.1 mJy for them. For MIPS 289, which is undetected, we assigned an upper limit of 0.1 mJy. We reject the source MIPS 8493 because of fringing in the VLA image at its position. The other four (MIPS 289, 15928, 16144, and 22530) are in regions where the radio images are clean. We correct their observed radio measurements to the rest frame, assuming a radio $K$-correction of $(1 + z)^{0.7}$, since a typical slope for a starburst radio spectrum is $-0.7$ (Condon et al. 1991). We plot the resulting $L(TIR)$ and $L(8 \mu m)$ values in Figure 8 as circles (filled for the detections, open for the upper limit.)

4.4.4. 8 $\mu m$/TIR Ratio for the $z \approx 2$ Lensed Galaxies

The lensed galaxies described in this paper bridge the luminosity divide between the $z \approx 2$ extremely high-luminosity galaxies in the literature, and the comparison LIRG and ULIRG samples at $z \approx 0$ and $z = 0.85$. For the lensed sample we can determine which galaxies have mid-IR spectra dominated by star formation, using the aromatic feature strengths. For the subset of lensed galaxies with submillimeter measurements and good upper limits at 70 $\mu m$, the SEDs are well enough constrained for a direct determination of $L(TIR)$. Therefore, the lensed galaxies can cleanly test for evolution with redshift in the behavior of the $L(8 \mu m)/L(TIR)$ ratio.

We use the aromatic-to-mid-IR flux ratio introduced in $\S$ 4.2 to select galaxies whose mid-IR spectra are dominated by star formation. The ratio selects A2219a, A1689b, A2218a, and A1835a as all having strong aromatic emission; their ratios range from $\approx 24 \times 10^{-11}$, with an average value of $16 \times 10^{-11}$. By comparison, our sample of local star-forming LIRGs and ULIRGs has a range in ratio of $\approx 7 \times 10^{-11}$ with an average of $21 \times 10^{-11}$. That is, the relative aromatic feature strength is virtually the same for the two samples, with no evidence that AGNs have significantly augmented the 8 $\mu m$ fluxes of our star-forming high-$z$ subsample. The EW(6.2 $\mu m$) criterion we used for the Yan et al. (2007) sample would have returned the same sample of star-forming lensed galaxies, except that A2219a lacks 6.2 $\mu m$ coverage and thus would not have been selected.

We plot the four lensed galaxies as red triangles in Figure 8.

4.5. Comparison of 8 $\mu m$/TIR Ratio at High and Low Redshift

The $L(8 \mu m)/L(TIR)$ ratio is important for two reasons: (1) it is one of the most accessible measures of the physical behavior of infrared emission as a function of redshift; and (2) many studies rely on fluxes at observed-frame 15 or 24 $\mu m$ to deduce $L(TIR)$ and related properties of galaxies at $z \approx 1$ or 2, and are thus relying on the behavior of $L(8 \mu m)/L(TIR)$. Figures 7 and 8 illustrate the important result that this ratio is well behaved for star-forming galaxies out to $z \approx 2.5$. In fact, this good behavior is a reasonably demanding test, since the various samples at $z \approx 2$ are subject to different selection effects and have $L(TIR)$ derived in different ways, yet they all follow a single trend. To recapitulate the different methods: the lensed galaxies are selected largely on amplification factor, and we measured $L(TIR)$ directly from template fits to rest frame mid and far-IR and submillimeter photometry. The four galaxies from Yan et al. (2007) were selected on the basis of colors and brightness at 24 $\mu m$, and we estimated $L(TIR)$ from the radio-infrared relation. The stacked galaxies from Papovich et al. (2007) were selected based on faint $K$-band and 24 $\mu m$ detections; we estimated $L(TIR)$ from rest 24 $\mu m$ measurements. The stacked galaxies from Daddi et al. (2005) were originally selected from BzK colors indicating they are at $z \approx 2$.

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20 See http://ssc.spitzer.caltech.edu/fls/extragal/vla.html.
and were stacked based on selections at a variety of wave-
lengths, and we used rest submillimeter measurements to esti-
mate $L_{\text{TIR}}$.

Having noted the overall uniform behavior, and the general simi-
larly of this behavior independent of redshift, we now eval-
uate the evidence from Figures 7 and 8 for modest evolution with
redshift. Specifically, there is a tendency for the ratio $L(8 \mu m)/
L_{\text{TIR}}$ to decrease with increasing luminosity for the local galax-
ies for $L > 10^{11} L_{\odot}$, but this tendency seems reduced for the
galaxies at $z \sim 2$. In fact, the high-redshift galaxies appear to have
SED behavior similar to normal local galaxies of lower luminosity
in this regard.

We have made two tests of the significance of this difference.
The first is based on the stacking results of Daddi et al. (2005)
and Papovich et al. (2007). Both studies select the galaxies to
stack on the basis of $K$-band detections, and therefore the mid-
and far-infrared characteristics of the galaxies should be represen-
tative of the entire population of IR-active galaxies (i.e., there is
no bias toward a specific $L(8 \mu m)/L_{\text{TIR}}$. To have an analogous
set of data for the local galaxies, we have averaged their $L(8 \mu m)/
L_{\text{TIR}}$ ratios in luminosity intervals of $10^{0.5}$; these averages are
shown as the large blue triangles in Figure 7. We have then fitted
the averaged data points, shown as a blue line. To test whether the
high-redshift stacked results were compatible with the same fit, we
carried out a second fit in which we included the three stacked
points. For the galaxies above $10^{11} L_{\odot}$, we then compared $\chi^2$ for the
low-redshift galaxies only in the two cases. We found that this $\chi^2$
is 4.5 times larger if the three stacked points at $z \sim 2$ are included in
the fit; that is, there is a strong indication that a single func-
tional fit cannot be used for both the local galaxies and for those
stacked at $z \sim 2$.21

To determine whether a few anomalous local galaxies could
affect this conclusion, we repeated the procedure but using a
more robust average for the luminosity-binned ratios. For each
bin with 10 or more items, we discarded the two highest and two
lowest ratios and then averaged the remainder. Where there were
less than 10 items, we retained the straight average. The high-
est luminosity bin, which is responsible for most of the change in
slope of the relation between the two luminosities, has 10 ob-
jects and hence its average was recomputed. The results were
virtually identical to those with the straight averages, as shown in
Figure 7.

A second test was applied to the high-redshift galaxies mea-
sured individually, either by Yan et al. (2007) or by us. In the case of
the Yan et al. sample, we expect a strong bias toward objects with
a large ratio of $L(8 \mu m)/L_{\text{TIR}}$, because they selected relatively
bright objects for spectroscopy, given the sensitivity lim-
itations of the IRS. Our sample was selected on a combination of
24 $\mu m$ flux density and lensing amplification, so a bias is likely,
although perhaps less severe.

To compare with these two studies, we selected a local sample
with relatively high 8 $\mu m$ output by removing all the galaxies that
are below the fit made to the averaged values (above). This sample
is shown in Figure 8. It is noteworthy that the scatter in $L(8 \mu m)/
L_{\text{TIR}}$ is small, with no outliers toward high $L(8 \mu m)$, so this is an
appropriate comparison sample for cases where a bias is suspected.
As before, we carried out two fits, one just to the local galaxies,
and another to them plus the eight individually measured high-
redshift ones. We found that $\chi^2$ for the low-redshift sample with
$L > 10^{11} L_{\odot}$ increased by a factor of 2.5 if we included the high-
redshift galaxies in the fit. Comparing Figure 8 with Figure 7, the
discrepancy would be substantially larger if we had included all the
local galaxies in our fits. This result confirms that found for the
stacked samples, namely, that a single functional fit is not a good
fit to the data over all redshifts.

Another comparison can be made with the sample of infrared
galaxies at $z \sim 0.85$ from Marcillac et al. (2006). The results fall
in a region occupied by both low- and high-redshift fits in Fig-
ure 8, which confirms the relatively modest change in $L(8 \mu m)/
L_{\text{TIR}}$ with redshift, but (not surprisingly given the small dif-
fERENCE between $z \sim 0$ and 2) does not definitively deviate from
either fit. In making this comparison, we have assumed that the
galaxies studied by Marcillac et al. (2006) have a bias toward
bright 15 $\mu m$ observed flux densities. If instead we treat them as
an unbiased sample and average them in luminosity intervals,
they fall along (or slightly above) the fit for the $z \sim 2$ sample in
Figure 7.

The tendency at high redshift for relatively strong 8 $\mu m$ emis-
ion could arise either through a contribution of an extra, AGN-
powered component at 8 $\mu m$ (e.g., the extra 5.8 $\mu m$ continuum
claimed by Sajina et al. 2007 and the mid-IR excess galaxies
with hard X-ray emission discussed by Daddi et al. 2007), or to a
change in the properties of the stellar-powered infrared-emitting
regions. The AGN explanation runs counter to our selection of
individual galaxies with strong aromatic features that correspond
to equivalent widths typical of purely star-forming local galaxies.
It is also counter to the precautions taken in the stacking analyses
to exclude AGNs (e.g., Papovich et al. 2007). In addition, AGN
contributions would come in a broad range, so it would be hard to
explain the consistent behavior of the five high-redshift samples in
Figures 7 and 8 if AGNs were causing the departure from the
behavior from local star-forming galaxies. Finally, the behavior
we see is similar to that found by Zheng et al. (2007) in a stacking
analysis at $z \sim 0.7$; specifically, they report that the far-IR SEDs
appear to resemble those of local galaxies of lower luminosity.
Thus, while AGNs may boost 8 $\mu m$ emission in other sources, it
is unlikely that AGN contamination is responsible for the behav-
ior seen in Figures 7 and 8.

We conclude that star-forming galaxies have a mild tend
toward increasing $L(8 \mu m)/L_{\text{TIR}}$ ratio with redshift that reaches
up to a factor of 2 at $z = 2$. In other words, $L(8 \mu m)$ does not
turn over with increasing $L_{\text{TIR}}$ for $z \sim 2$ galaxies as it does at
low redshift. This probably helps explain why some high-redshift
galaxies have higher 6.2 $\mu m$ aromatic equivalent widths than
are observed for local ULIRGs (Desai et al. 2007; Menéndez-
Delmestre et al. 2007; Sajina et al. 2007), although direct com-
parisons cannot be made because there are not local galaxies with
such extremely high luminosities. Because the lensed galaxies
and stacked samples directly probe the $10^{11} - 10^{12.5} L_{\odot}$ range at
$z \sim 2$, we can directly compare to $z \sim 0$ galaxies at similar lumin-
osity rather than extrapolating.

That aromatic luminosities do not turn over with $L_{\text{TIR}}$ in
high-redshift galaxies may be related to lower densities and op-
tical depths of the infrared-emitting regions in the high-$z$ galax-
ies. Lower optical depths would result in part if the metallicity
were lower, resulting in a higher gas-to-dust ratio. Draine et al.
(2007) show that the dust content of galaxies has a very steep
dependence on their metallicity. The densities and optical depths
of the star-forming regions would also be lower if the high lumin-
osities are associated with star formation on large spatial scales
in the high-redshift galaxies, as opposed to strongly concentrated,
major-merge—associated nuclear star formation in local ULIRGs.

21 This analysis makes the assumption that low-luminosity galaxies have
similar behavior at low and high redshift.
Such very dense regions in local ULIRGs appear to account for the scattering of points below the \( z = 0 \) trend line. Use of local templates to estimate \( L(\text{TIR}) \) for high-\( z \) galaxies detected at 24 \( \mu \)m has indicated the existence of very high luminosities at \( z \approx 2 \) and beyond (e.g., Yan et al. 2007). The change in \( L(8 \mu m)/L(\text{TIR}) \) suggested here would reduce \( L(\text{TIR}) \) for a given \( L(8 \mu m) \) by a factor of \( \sim 4 \). As shown in the figure, in our sample there is only one galaxy selected to be dominated by star formation that has a \( L(\text{TIR}) \) significantly exceeding \( 10^{13} L_\odot \). A luminosity of \( 10^{13} L_\odot \) corresponds to a star formation rate of \( \sim 2000 M_\odot \) \( \text{yr}^{-1} \) if we use the local calibration (Kennicutt 1998), and is suggestively close to the theoretical maximum luminosity that a starburst is allowed by causality arguments—that star formation throughout a galaxy should not transpire faster than the dynamical timescale (Lehnert & Heckman 1996; Elmegreen 1999). (The actual maximum value depends on efficiency and dynamical mass.) Thus, our lower calibration for \( L(8 \mu m)/L(\text{TIR}) \) may remove the discrepancy between \( L(\text{TIR}) \) inferred for high-redshift galaxies and plausible upper limits on what a starburst can produce. Spatially resolved imaging of high-redshift starbursts (e.g., Nesvadba et al. 2007) can test whether they are in fact forming stars at the dynamical limit.

5. SUMMARY AND IMPLICATIONS

The 24 \( \mu \)m photometric band of Spitzer efficiently detects galaxies out to \( z \approx 3 \), and is being used as an extinction-robust measure of star formation rates in the distant universe. We present new IRS spectra and MIPS photometry of gravitationally amplified galaxies, and use them in concert with archival Chandra images and published submillimeter photometry to investigate the nature of the very faint 24 \( \mu \)m sources.

The aromatic feature flux ratios observed in our spectra agree well with those of star-forming template spectra at \( z \approx 0 \). This result implies that dust size and ionization distribution are not strongly evolving with redshift. Thus, for \( z \approx 1 \) star-forming galaxies with little AGN contribution, the mid-IR spectra can be approximated by \( z \approx 0 \) star-forming templates.

22 This galaxy was selected as star-forming based on its EW(6.2 \( \mu \)m) (Saikina et al. 2007), but its EW(7.7 \( \mu \)m) is unusually low, which may indicate a substantial AGN contribution.

At \( z \approx 2 \), the 24 \( \mu \)m data probe the 8 \( \mu \)m rest wavelength region. We find that the ratio of rest-frame 8 \( \mu \)m luminosity to total infrared luminosity is approximately the same as for local galaxies of similar \( L(\text{TIR}) \). However, the \( z \approx 0 \) relationship does not appear to hold perfectly. Combining data from the literature with this sample, we demonstrate that \( z \approx 2 \) star-forming galaxies have higher \( L(8 \mu m) \) luminosities for a given \( L(\text{TIR}) \) than low-redshift analogs. Functionally, the break in the \( L(8 \mu m) \) versus \( L(\text{TIR}) \) relationship occurs at a 10-times-higher \( L(\text{TIR}) \) at \( z \approx 2 \) than it does at \( z \approx 0 \). We speculate that this is due to lower optical depths in the IR-emitting regions of the high-\( z \) galaxies, perhaps because of lower metallicity or because star formation is more spatially extended than in local ULIRGs (e.g., across a disk rather than in a compact nucleus). This result has important implications for mid-IR--derived star formation rates at \( z \approx 2 \), and may indicate a maximum star-forming luminosity of a few \( \times 10^{13} L_\odot \).

The results just summarized were found for sources where AGNs do not contribute substantially to the mid-IR spectra, as determined by the aromatic-to-mid-IR flux ratio. But 6 of the 19 mid-IR spectra in our extended sample do show a strong AGN. Including other AGN diagnostics (X-ray detections and optical emission lines), 8 of the 19 galaxies show evidence for nuclear activity. The prevalence of AGNs in this sample may contradict the frequent assumption that the 24 \( \mu \)m flux density in 100 \( \mu \)Jy sources at \( 1 < z < 3 \) is dominated by star formation, with AGNs a modest contaminant.

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APPENDIX

NOTES ON INDIVIDUAL SOURCES

A1. SOURCE A2390a

J. Richard obtained a Keck spectrum for this source, shown in Figure 9. The strong Ly\( \alpha \) emission is redshifted by \( z = 2.858 \), and the line velocity spread (FWHM \( \sim 945 \text{ km s}^{-1} \)) indicates that the source is an AGN.

A2. SOURCE A2261a

This submillimeter galaxy (Chapman et al. 2002) lacks a spectroscopic redshift; submillimeter and radio-based photometric estimates place it at \( z > 1.6 \)(Chapman et al. 2002) and \( 2.5 < z < 4 \)(Arexaga et al. 2003). It is the only source without a redshift in our sample. Our spectrum should be sensitive to the 11 \( \mu \)m complex for \( 0.9 < z < 1.8 \), and to the 7.7 \( \mu \)m aromatic feature for \( 1.8 < z < 3.1 \). Thus, for the aromatic features to fall outside the IRS spectral coverage, the source would have to have very low \( z < 1 \) or very high \( z > 3.1 \) redshift. If the source were at \( z > 3.1 \), it must be \( \sim 1000 \) times more luminous than M82 to have produced the observed 24 \( \mu \)m flux density. It seems more probable that this source lies at \( 1 < z < 3 \) and lacks aromatic emission.
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Fig. 9.—Keck/LRIS spectrum of A2390a. The strong feature at $\lambda = 4690$ Å is presumably Ly$_\alpha$ at $z = 2.858$. 

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