The proper motion of PSR J0205+6449 in 3C 58

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1 INTRODUCTION

The supernova remnant 3C 58 (G130.7+3.1) has a filled centre morphology and a flat radio spectrum, and was classified on this basis as being pulsar wind nebula (PWN, also plerion; e.g., Weiler & Panagia 1978) long before a pulsar was seen. The pulsar, PSR J0205+6449, was subsequently detected, first in the X-ray (Murray et al. 2002) and then in the radio (Camilo et al. 2002). It is also one of the approximately 100 pulsars that shows pulsed gamma-ray emission (Abdo et al. 2009).

PSR J0205+6449 and 3C 58 are at a distance, D, of ~3.2 kpc (Roberts et al. 1993).1 They have traditionally been associated with the historical supernova of 1181 AD (SN 1181; Clark & Stephenson 1977; Stephenson & Green 1999), giving them an age of ~830 yr and making them one of the youngest known supernova remnants and pulsars. However, recent work has suggested that 3C 58 is likely considerably older. In particular, the measured expansion speeds of both the synchrotron bubble (Bietenholz, Kassim & Weiler 2001; Bietenholz 2006) and of the thermal filaments (Fesen 1983; Fesen, Kirshner & Becker 1988; van den Bergh 1990; Rudie & Fesen 2007), seem to be considerably lower than expected for an age of ~830 yr, suggesting an age of several thousand years. An age several times larger than 830 yr is also suggested by several other arguments (see e.g. Chevalier 2005; Bietenholz 2006).

PSR J0205+6449 has a spin frequency of ~15.2 Hz, with a derivative of ~−4.5 × 10−11 Hz s−1 (Camilo et al. 2002; Murray et al. 2002; Livingstone et al. 2009), so the characteristic age of the pulsar is 5400 yr. (We note that if the pulsar had an initial spin frequency of ~16.3 Hz, the present spin frequency and spin-down rate could be reconciled with an age of 830 yr; however, as mentioned, numerous other arguments independent of the characteristic age suggest a considerably larger age.) It has a high level of timing dispersion.
noise, having exhibited two spin-up glitches with fractional magnitudes of $\Delta v/v = 3.4 \times 10^{-7}$ and $3.8 \times 10^{-6}$ between MJDs 52276 and 53063 (Livingstone et al. 2009). The pulse width is $\sim 2.5$ ms (at 1.4 GHz; Camilo et al. 2002). Although as mentioned, there is some controversy over the age of 3C 58, PSR J0205+6449 is nonetheless one of the youngest known pulsars. 

PSR J0205+6449 is a particularly interesting case because of the presence of the easily observable PWN, 3C 58. Such nebulae, seen for only a handful of pulsars, provide important diagnostics for young pulsars, giving insight into the winds which carry away the pulsar spin-down energies.

Knowing PSR J0205+6449’s proper motion is important for determining its exact birthplace as well as the nature of its interaction with its surrounding PWN. More generally, the origin of the high space velocities of pulsars is an interesting question, and therefore knowledge of PSR J0205+6449’s proper motion is particularly important because of its young age.

2 PULSAR TIMING WITH GBT: OBSERVATIONS AND RESULTS

We obtained two sessions of pulsar timing observations at the $\sim 105$ m diameter Robert C. Byrd Green Bank Telescope (GBT), which were carried out simultaneously with the very long baseline interferometry (VLBI) sessions. We used a new observing technique where the data from the GBT were simultaneously used for both pulsar timing and the VLBI observations. The VLBI data are described in the next section.

On 2007 April 25, we observed PSR J0205+6449 with the GBT Spigot backend (Kaplan et al. 2005) for a duration of 10.8 h. The Spigot measured three-level autocorrelations from each of two polarizations for 1024 lags and integrated them for 81.92 s covering 800 MHz, centred at 1650 MHz, approximately 500 MHz of which (from $\sim$1350 to 1850 MHz) were free enough from interference to be used for the timing analysis. We let the Spigot continue to take data while the GBT occasionally moved to and from a calibrator source as part of the VLBI scheduling (see below). For these times, which correspond to about 35 per cent of the total duration, we simply zero-weighted the resulting data from the pulsar. A standard pulsar timing analysis using TEMPO provided us with a pulsar spin ephemeris, which we then used for pulsar gating of the interferometric data during correlation. We measured a barycentric pulsar spin period of 0.0657167116(4) s at epoch MJD 54215.3220 (UTC).

We performed a very similar observation with the pulsar backend Green Bank Ultimate Pulsar Processing Instrument (GUPPI; Ransom et al. 2009) on 2010 October 18, where we observed PSR J0205+6449 for 8.9 h. GUPPI recorded the 8 bit, summed polarizations from 2048 channels every 192 $\mu$s covering 800 MHz of bandwidth (of which the highest $\sim$600 MHz was usable). After zero-weighting data when the GBT was off of the pulsar’s position, a timing analysis determined the barycentric pulsar spin period to be 0.0657388555(6) s at epoch MJD 55487.8302 (UTC).

For our timing analyses we used a simple two-Gaussian model as our template profile where the peak of the pulsar flux occurred very near to spin phase 0.5. We assumed a pulsar position, dispersion measure and instantaneous spin-down rates consistent with those measured by Livingstone et al. (2009).

We show the average pulse profiles obtained in the two observing sessions in Fig. 1. As the pulsar is quite weak, the signal-to-noise ratio is relatively low, and consequently we do not consider the differences between the two profiles significant. The full width at half-maximum (FWHM) duty cycle of the pulse is $\sim 6$ per cent.

3 VLBI RADIO OBSERVATIONS

We obtained two epochs of 1.4-GHz VLBI observations of PSR J0205+6449 using the ‘High Sensitivity Array’, which consisted of the NRAO Very Long Baseline Array (VLBA) augmented by the GBT ($\sim 105$ m diameter) and Effelsberg (100 m diameter) telescopes. For the first session, on 2007 April 25 (program code BB241), we obtained 10 h of VLBI data, while for the second session on 2010 October 18 (program code BB295), we obtained 12 h. In each session we recorded both senses of circular polarization, and we used 2-bit sampling at a bit rate of 512 Mbit s$^{-1}$, for an effective bandwidth per polarization of 64 MHz.

The data from 2007 were correlated with NRAO’s VLBA processor, while those from 2010 were correlated with the ti$\beta$x processor (Deller et al. 2011). The analysis was carried out with NRAO’s Astronomical Image Processing System (AIPS). During correlation, we gated the correlator using the pulsar timing extracted from the GBT.

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2 Several pulsars or PWNe, including the Crab nebula, G21.5−0.9 (Bietenholz & Bartel 2008) and Kes 75 (Gotthelf et al. 2000) are thought to be only around 1000 yr old, but the vast majority of known pulsars have ages $>10 000$ yr.
data (as described above), and accumulated several sets of VLBI data with different gating: first, the ungated data, and secondly, the ‘on-pulse’ gated data where the gated bin had a width of 8 per cent of the pulsar period (distributed symmetrical about the peak). In addition, for the 2010 data, correlated with the TECOR correlator, we also accumulated data in 12 different bins evenly distributed in pulse phase, with each bin therefore having a width of 8.33 per cent of the period. In all data sets we used a correlator integration time of 3.0 s.

We used VCS2 J0209+6437 (Fomalont et al. 2003), 0.5 away from PSR J0205+6449, as our primary phase-reference source, but we included some observations of an astrometric check source, JVAS J0228+6721(4C 67.05, ICRF J022850.0+672103; IERS B0224+671), which is a source in the International Celestial Reference Frame (ICRF) catalogue whose position is accurately known (Fey, Gordon & Jacobs 2009), and which is 3.4 away from PSR J0205+6449. For the phase-centre position for the pulsar observations, we use the position of the pulsar given in Slane, Helfand & Murray (2002), 02°50′37″92, 64°49′42″8. Mostly we used a cycle time of 3.8 min, with 2.5 min being spent on PSR J0205+6449. The calibration of the VLBI data was done using standard procedures using NRAO’s AIPS package. The flux density calibration was done through measurements of the system temperature at each telescope, and the antenna amplitude gains were subsequently improved through self-calibration of the reference sources.

Some sporadic Radio Frequency Interference (RFI) was seen, particularly in the frequency range 1617–1625 MHz. We clipped out visibility points with anomalously high amplitudes, removing \( \lesssim 1 \) per cent of the data. We also deleted any data taken when either of the antennas were at elevations of \( \leq 10° \). During our 2007 observing run, the GBT failed to observe J0209+6437 for three periods of approximately 1 h. We discarded the GBT data for all sources from these periods.

The fringe fitting and phase calibration were done using the (ungated) data from our calibrator sources. We estimated the ionospheric delay using the AIPS task TECOR using IONEX data from the Crustal Dynamics Data Information System archive3 of the Goddard Space Flight Center. The earth orientation parameters used during the correlation were extrapolated. We corrected to more accurate ones subsequently made available from the United States Naval Observatory. The resulting calibration was then interpolated to the times of the PSR J0205+6449 as well as the J0228+6721 observations, and applied to both the gated and ungated visibility data for PSR J0205+6449.

3.1 Astrometric reference sources

Our primary phase-reference source was J0209+6437. We take, as the astrometric reference point for J0209+6437, the peak brightness position, for which we take the coordinates 02°50′35°98806(20), 64°37′25′′7701(06) from the VLBA Calibrator Survey 2 (VCS2; Fomalont et al. 2003), where the numbers in parentheses give the uncertainty in the last two digits. Unfortunately, this source turned out to be somewhat resolved. We show the VLBI images of J0209+6437 on our two observing epochs in Fig. 2. The source is elongated in the north–south direction, with the 50 per cent contour having an extent of approximately 6 mas.

It is possible that the peak brightness position, and therefore our astrometric reference point, varies with time due to changes in the source morphology (see e.g. Bietenholz, Bartel & Rupen 2000; Bartel et al. 2012).

As mentioned, we also included some observations of an astrometric check source, J0228+6721. Like our observations of PSR J0205+6449, we phase reference our observations of J0228+6721 to J0209+6437. We give the peak brightness positions of J0228+6721, as well as the ICRF position in Table 1.

Unfortunately, this source is also extended at the mas level in the north–south direction (Perley, Fomalont & Johnston 1980; Romney et al. 1984; Lister et al. 2009a). Moreover, this source is known to have a somewhat variable morphology (Padrielli et al. 1986; Bondi et al. 1996; Lister et al. 2009b), with both changes in the size of the central component, and motions of other components being observed with magnitudes of up to \( \sim 0.3 \) mas yr\(^{-1} \). We show images of J0228+6721 in Fig. 3.

3.2 PSR J0205+6449

The pulsar was expected to be quite faint: even in the images made from gated VLBI data, we expected a signal-to-noise ratio of \( \lesssim 10 \). Furthermore, the pulsar position is not accurately known a priori. In order not to bias our analysis by searching only near an assumed position, we therefore chose to image a fairly large region. We imaged PSR J0205+6449 directly from the calibrated multichannel VLBI data without any averaging of the visibility data in frequency. Since our channel width was 0.5 MHz, the field of view over which bandwidth smearing is small is \( \sim 10 \) arcsec, which is adequate for our purposes. We made images of 4096 \times 4096 pixels, with the pixels being 0.8 \times 0.8 mas, centred on our initial position estimate for PSR J0205+6449, which was the position used in the correlation, and was the one given by Slane et al. (2002) of 02°50′37′′92, 64°49′42′′8.

\[ \text{http://cddis.gsfc.nasa.gov} \]
Table 1. Position of J0228+6721.

| Observations         | Frequency (GHz) | RA (J2000) (h m s) | Dec. (J2000) (° ′ ″) | Offset (mas) |
|----------------------|-----------------|--------------------|----------------------|--------------|
| ICRF (Fey et al. 2009) | 8.4             | 02 28 50.05148948  | 67 21 03.0293039     |              |
| 2007 VLBI obs. a      | 1.4             | 02 28 50.0516247   | 67 21 03.03364       | 0.78          |
| 2010 VLBI obs. a      | 1.4             | 02 28 50.0516134   | 67 21 03.05272       | 0.72          |

aOffset of measured position from the ICRF position.

bPositions measured from peak brightness point, and measured relative to J0209+6437, which is assumed to be at the VCS2 position (see text).

Table 2. Position of PSR J0205+6449.

| Observation epoch | Gating | Image extrema (in units of image rms) | RA (J2000) (h m s) | Dec. (J2000) (° ′ ″) |
|-------------------|--------|--------------------------------------|--------------------|----------------------|
| 2007              | Off-pulse | −5.1 +5.1                          | (Not detected)     | 64 49 41.3319 (4)   |
| 2007              | On-pulse  | −5.0 +6.8                           | 02 05 37.92322 (1) | 64 49 41.3319 (4)   |
| 2010              | Off-pulse  | −5.5 +5.2                          | (Not detected)     | 64 49 41.3343 (4)   |
| 2010              | Off-pulse  | −5.1 +6.5                           | 02 05 37.92247 (1) | 64 49 41.3343 (4)   |

The position of the brightness peak of PSR J0205+6449. The digit in parentheses gives the statistical uncertainty in the last digit of the corresponding coordinate value; note that the total uncertainty is dominated by a systematic component discussed in the text.

In Table 2 we give the signal-to-noise ratios of the brightness extrema in the various images. In the images made from ungated visibility data, the positive and negative extrema were of similar magnitude, being ≤5.5σ, where σ was the image brightness rms. In the images made from the on-pulse gated data, for both epochs, the negative extrema were of a similar magnitude. For both epochs, however, the images made from the on-pulse data showed a single positive peak that was ≥6.5σ, or larger in magnitude than the negative extremum by ≥1σ. The number of independent sky positions sampled by our images can be estimated by dividing the image area by the FWHM area of the fitted beam and was ∼650 000. If the pixel brightnesses are Gaussian distributed, the likelihood of obtaining a value in excess of ±6σ value in 650 000 trials is then ∼0.1 per cent, although we note that the distribution of pixel brightnesses may not be accurately Gaussian at these large deviations from 0.

Furthermore, the two positive extrema seen in the images made from the gated data were quite close to one another on the sky, being within 6 mas of each other on the sky. Since the total area examined was ∼107 mas2, the odds of random peaks occurring so close together would be ∼80 000 to 1 (which probability is independent of the actual distribution of pixel brightnesses).

In Fig. 4 we show a section of the images made from the gated on-pulse and ungated visibility data sets from 2010. Although the signal-to-noise ratio is not high, the pulsar can be clearly seen in the on-pulse image, but not in the corresponding ungated one. Despite the modest signal-to-noise ratio, we think that the detections are firm for several reasons. First, in each of the ‘on-pulse’ images, the positive extremum in the image (excluding a ∼100-pixel strip around the edge where the rms is somewhat higher due to numerical artefacts) is ≥1σ higher than the corresponding negative one, which is not expected to happen by chance. Secondly, no source is visible in either of the ungated images at this location. Any background source should be much more strongly detected in the images made from the ungated data since the gated visibility data represent only ∼8 per cent of the total observations.

As an additional check, for the 2010 observations, we imaged separately the visibility data from the 12 bins spaced...
Figure 4. VLBI images of PSR J0205+6449 from our 2007 observations. Both images are “dirty” or underconvolved. In both panels, the grey-scale runs from $-3.8\sigma$ (white) to $6.9\sigma$ (black), and the FWHM of an elliptical Gaussian fit to the central part of the beam is shown at lower left. The left-hand panel shows the image made from the on-pulse gated visibility data, while the right-hand panel shows the image made from the ungated visibility data. The images are centred on PSR J0205+6449, only visible in the left-hand panel, which is at $02^h05^m37.92322, 64^\circ49'41.3319$.

Figure 5. The pulse profile of PSR J0205+6449: the image brightness, given in units of the image rms ($\sigma$), at the location of PSR J0205+6449 in the images made from the data in each of 12 bins distributed across the pulsar period. The first bin on the left is duplicated on the right for clarity. The shaded area shows the $\pm1\sigma$ region. Only in the bin at pulse phase $= 0.5$ is a signal substantially in excess of the noise visible.

3.3 Proper motion of PSR J0205+6449

Using the position determinations in Table 2 we obtained the displacement of PSR J0205+6449 between our two observing epochs to be $-4800 \pm 570$ and $+2800 \pm 570$ $\mu$as in RA and Dec., respectively, where the uncertainties are statistical only. At a distance of 3.2 kpc, the parallax of PSR J0205+6449 is expected to be 313 $\mu$as, smaller than the statistical uncertainty, but as the angular displacement due to parallax is readily calculable we corrected for it. The expected shift between our two particular epochs due to parallax is to be 36 and 499 $\mu$as in RA and Dec., respectively. Correcting for the parallax, we then arrive at a net displacement of PSR J0205+6449 between our two epochs of $-4860 \pm 570$ and $+2330 \pm 570$ $\mu$as in RA and Dec., respectively (where we have ignored the small additional uncertainty incurred due to the uncertainty in the distance). This displacement corresponds to a proper motion of $-1400 \pm 160$ $\mu$as yr$^{-1}$ in RA and $540 \pm 160$ $\mu$as yr$^{-1}$ in Dec., or $1500 \pm 160$ $\mu$as yr$^{-1}$ at position angle (PA) $159^\circ \pm 6^\circ$, where the proper motion is measured relative to the brightness peak of J0209+6437, and the uncertainties are statistical only.

Since our reference source, J0209+6437, is somewhat resolved, we must carefully assess any possible effect of temporal changes in J0209+6437 on our proper motion measurements. In our VLBI observations, we observed a check source, J0228+6721 in a fashion similar to PSR J0205+6449, in particular similarly phase referenced to J0209+6437. The positions measured for the check source were given in Table 1, and suggest a proper motion for J0228+6721 of $-1$ and $550$ $\mu$as yr$^{-1}$ in RA and Dec., respectively. The source J0228+6721 is an ICRF source and Feissel-Vernier (2003) find a proper motion of $\leq 50$ $\mu$as yr$^{-1}$ at 8.4 GHz over a period of $\sim 12$ yr. However, the lack of secular motions at 8.4 GHz, where the core is likely more dominant, does not preclude apparent motions at our lower frequency of 1.4 GHz, where jet components are more uniformly across the pulsar period. In Fig. 5 we show brightness at the location determined from the on-pulse image as a function of the pulse phase. The profile can be compared to that obtained from the GBT pulsar-timing observations shown in Fig. 1.

We note here that position of PSR J0205+6449 is different than that of the point source visible in the X-ray as given in Slane et al. (2002) of $02^h05^m37:92, 64^\circ49'42.8$ by $\sim 1.5$ arcsec.
likely to dominate the emission. Indeed, other authors have reported larger proper motions for components within J0228+6721: Bondi et al. (1996) find an expansion of 300 µas yr⁻¹ while Lister et al. (2009b) find several moving components with proper motions of up to 376 µas yr⁻¹ at 14 GHz. It is therefore possible that our observed proper motion is due to component motions within J0228+6721. It is, however, also possible that our primary reference source, J0209+6437, is not in fact stable. Unfortunately, both our reference sources are elongated in the north–south direction, so that the observed north–south relative motion between them could be the result of motions along the jet axis in either source.

As a conservative estimate of 1σ uncertainty on the proper motion of PSR J0205+6449, as phase referenced to J0209+6437, we therefore take the apparent proper motion of J0228+6721 found above. We then arrive at a final value for the proper motion of PSR J0205+6449 of −1400 ± 160 µas yr⁻¹ in RA and 540 ± 575 µas yr⁻¹ in Dec. At a distance of 3.2 kpc, this corresponds to a speed of 23 ± 6 km s⁻¹. Note, however, that this velocity is with respect to the Earth. We calculate a more physical value of PSR J0205+6449’s projected velocity, corrected for Galactic rotation and the Sun’s motion, in Section 5.

4 X-RAY ASTROMETRY

To determine the X-ray position of PSR J0205+6449, we investigated data from a deep Chandra observation carried out between 2003 April 22 and 2003 April 26 (see Slane et al. 2004, for details and original results from these observations). The data from observation IDs 4383, 4382 and 3832 were reprocessed and cleaned using standard routines from CIAO version 4.4. The merge_all task was used to create a merged image from the 317 ks of good exposure time. The pulsar is embedded in a bright compact nebula that is slightly asymmetric in the east–west direction. The centroid of the point source emission is located at a position of 02¹⁰⁰⁵³⁷.93, +64°49′41″.4. We identified three X-ray point sources in the field that have counterparts in the Two Micron All Sky Survey (2MASS) catalogue. Comparing the centroid positions of the X-ray sources with the infrared positions, we find that the uncertainty in the Chandra position of the pulsar is σ_RA ≈ 0.06 s and σ_Dec. ≈ 0.12 arcsec. The systematic error associated with contributions from the compact nebula is estimated to be slightly larger than these values, but the combined uncertainty is ≲ 1 arcsec either direction.

We note that position given above is slightly different from that determined in the shorter Chandra observation reported by Slane et al. (2002). We have reprocessed those earlier data (observation ID 728) as well, and find that the position is in excellent agreement with the value given above. We conclude that there were small errors in the initial position reconstructions reported by Slane et al. (2002).

5 DISCUSSION

Using gated VLBI we imaged PSR J0205+6449 in the centre of 3C 58. The pulsar was detected. The position of PSR J0205+6449 we found from the VLBI observations was 1.5 arcsec distant from the originally published one of a compact X-ray source seen in Chandra Advanced CCD Imaging Spectrometer (ACIS) observations (Slane et al. 2002). A re-examination of the Chandra data, however, resulted in an improved position of the X-ray source which is well within the uncertainties of that measured with VLBI.

Shearer & Neustroev (2008) report deep optical observations of the centre of 3C 58 using the 4.2-m William Herschel Telescope in La Palma, in which they detect three unresolved sources with R-band magnitudes of ~24, in addition to the more diffuse synchrotron emission from the PWN. Shearer & Neustroev (2008) tentatively identified their source ‘o1’ as PSR J0205+6449 on the basis of its coincidence with the published position of the compact X-ray source. However, based on our VLBI determination and the re-determined Chandra High Resolution Camera (HRC) position, it is in fact their object ‘o2’, which is almost certainly the optical counterpart of PSR J0205+6449. The position of o2 is RA = 02°05′37.93 and Dec. = +64°49′41″4, with an uncertainty of ≲ 0.1 arcsec, consistent to within 0.08 arcsec with the pulsar position determined from VLBI. Shearer & Neustroev (2008) give the magnitudes of o2 as 24.15 ± 0.07 in R, >24.3 in V and >25.6 in B, consistent with the optical magnitudes estimated for PSR J0205+6449 based on its spin-down luminosity, distance and the expected Galactic extinction.

In summary, we can regard the detection of PSR J0205+6449 in the radio, optical and X-ray bands as firm, with the most accurate position being that determined from the VLBI observations and given in Table 2.

In addition to determining the position, we also determined the proper motion of PSR J0205+6449 of −1400 ± 160 µas yr⁻¹ in RA and 670 ± 575 µas yr⁻¹ in Dec. In order to correct this measured value for Galactic rotation and the Earth’s motion with respect to the local standard of rest, we take the Galactic constants from Schönrich (2012), namely a flat Galactic rotation curve with v = 238 km s⁻¹ and a distance to the Galactic Centre of 8.27 kpc. We take also from Schönrich (2012) the solar motion of 250 km s⁻¹ in the Galactic plane, solar radial velocity of 14 km s⁻¹ towards the Galactic Centre and a motion of 6.1 km s⁻¹ perpendicular to the Galactic plane. We take the distance of PSR J0205+6449 (from the Earth) to be 3.2 kpc. We can then calculate the peculiar motion of PSR J0205+6449 with respect to the standard of rest at its location as being 4° 2.3 ± 0.3 mas yr⁻¹, corresponding to 35 ± 6 km s⁻¹ at PA = −38° (PA in equatorial rather than Galactic frame).

This speed is relatively small: Galactic pulsars have a two-dimensional velocity dispersion of 200–300 km s⁻¹ (e.g. Lyne & Graham-Smith 1990; Hobbs et al. 2005). Even taking the uncertainty introduced by possible motion of our reference source, the 3σ upper limit on the proper motion of PSR J0205+6449 is only 53 km s⁻¹.

Can PSR J0205+6449’s low tangential velocity be reconciled with the velocity distribution of other Galactic pulsars? Faucher-Giguère & Kaspi (2006) found that the distribution of space velocities of Galactic pulsars at birth was consistent with one where each of the three orthogonal components of the space velocity was distributed exponentially, with the mean of the absolute value of each component being 180 ± 20 km s⁻¹, and the resulting mean three-dimensional speed being 380 ± 40 km s⁻¹. Since PSR J0205+6449 is quite young (we discuss the age in more detail below), its velocity is probably close to its birth velocity. As the multidimensional exponential distribution of Faucher-Giguère & Kaspi (2006) is difficult to evaluate numerically, we performed a Monte Carlo simulation with n = 10 000 trials, to find that the chance of finding a pulsar with a tangential velocity as small as the 35 ± 6 km s⁻¹ measured for PSR J0205+6449 from such a distribution is ~2.7 per cent.

4 We noted above that the proper motion uncertainty is rather larger in Dec. than in RA. Here we take the geometric mean of the two values.
Figure 6. A 1.4 GHz VLA radio image of 3C 58, reproduced from Bietenholz (2006). The ‘×’ sign shows the present location of PSR J0205+6449, while the two ‘+’ signs show the extrapolated position at two possible dates of the supernova explosion which gave rise to 3C 58: at 1181 AD and at 5000 BC, with the one at 1181 AD being the one overlapping with the present position (×). The ‘○’ shows the position of the centre of the circular region of softer X-ray emission, interpreted as thermal emission from the supernova shell, from Gotthelf, Helfand & Newburgh (2007). The grey-scale is labelled in mJy bm⁻¹, and the FWHM resolution was 1.4 arcsec. We show a detail of the central region in Fig. 7.

Alternatively, Hobbs et al. (2005) collected proper motion measurements for 140 pulsars. Of their sample, 8 or 5.7 per cent had tangential velocities ≤35 km s⁻¹. Finally, some authors have found the distribution of pulsar velocities to be bimodal. For example, Brisken et al. (2003) find that 20 per cent of pulsars form a low-velocity component with a one-dimensional velocity dispersion, σv1D of 99 km s⁻¹, while the remainder have a σv1D = 294 km s⁻¹. In this case, the probability of a random pulsar having a tangential velocity as low as 35 km s⁻¹ is 1.8 per cent.

In conclusion, the small measured proper motion makes PSR J0205+6449 somewhat unusual in being amongst the slowest few per cent of young pulsars regardless of whether a bimodal or a unimodal distribution of pulsar velocities is considered.

Our estimate of the tangential velocity depends on the distance to PSR J0205+6449 and 3C 58, which, as mentioned earlier, is somewhat uncertain. The value of 3.2 kpc of Roberts et al. (1993) was determined kinematically from H I absorption. Kothes (2010) argues for a distance of ~2 kpc. Adopting a smaller distance would reduce the tangential velocity estimate by the corresponding factor, and thus make PSR J0205+6449’s low speed even more unlikely.

The low measured angular speed therefore argues against a distance much lower than 3.2 kpc.

On the other hand, PSR J0205+6449’s dispersion measure is approximately twice that expected for a distance of 3.2 kpc (Camilo et al. 2002), suggesting perhaps that the true distance is somewhat larger, which would imply a tangential velocity larger by a factor of (D/3.2 kpc). This question could be resolved with a measurement of PSR J0205+6449’s trigonometric parallax with 10 per cent accuracy. Such a measurement is feasible if an in-beam calibrator can be found (see e.g. Chatterjee et al. 2009), and should be undertaken.

Can we extrapolate back from PSR J0205+6449’s present position to determine its position at the time of the supernova explosion? To do so requires knowing the age of PSR J0205+6449. As noted in the Introduction, although it has traditionally been associated with a supernova in 1181 AD, making it ~830 yr old, it is probably older, with a likely age of ~7000 yr (Bietenholz et al. 2001; Bietenholz 2006), suggesting an epoch of ~5000 BC for the supernova event. In Fig. 6 we show a radio image of 3C 58, and indicate the present, and extrapolated epoch 1181 AD and 5000 BC positions of PSR J0205+6449. Fig. 7 is a detail of only the central region of Fig. 6, showing the pulsar positions more clearly.

From XMM–Newton X-ray observations of 3C 58, Gotthelf et al. (2007; see also Bocchino et al. 2001) found an approximately circular region with a softer X-ray spectrum, which they interpreted as a thermal X-ray emission from the supernova shell. They noted, however, that the centre of this shell was displaced from the pulsar...
position at 02°55′33″97, 63°49′50″0 (J2000.0). We plot this centre position also in Fig. 6. Neither the present distribution of synchrotron emission nor thermal filaments suggests any particular explosion centre, and although Fesen et al. (2008) determined proper motions for various optical features, the precision is too low to accurately identify an expansion centre.

Our measured proper motion does not place PSR J0205+6449 near the centre of the region of softer X-ray emission for either possible explosion epoch. For an explosion epoch of 5000 yr, however, the extrapolated position of PSR J0205+6449 is closer to the present geometrical centre of 3C 58.

If the area of softer X-ray emission identified by Gotthelf et al. (2007) were thermal X-ray emission associated with the forward shock of the supernova, then one would probably expect its centre to be near the location of the explosion and of PSR J0205+6449’s birth. This does not seem to be the case. Furthermore, the diameter of the region of softer X-ray emission is only \( \sim 5.6 \) pc, which is notably smaller than the PWN, which has an east–west extent of \( \sim 8.5 \) pc. The PWN, however, is expected to still be confined by the ejecta and thus be inside the forward shock, which suggests that the forward shock is considerably larger than the region of softer X-ray emission. Therefore, both because the region of softer X-ray emission is not centred on the location of the explosion, and because it is small compared to the PWN, we think it is unlikely that the softer X-ray emission is associated with the supernova forward shock.

6 SUMMARY AND CONCLUSION

(1) We obtained VLBI observations of PSR J0205+6449, the pulsar in 3C 58. We employed a novel technique to obtain VLBI observations of faint pulsars, where we used the GBT simultaneously for pulsar timing observations as well as an element of the VLBI array. The derived pulsar timing information was then used to gate the VLBI correlator increasing the signal-to-noise ratio of the pulsar VLBI. This technique can be used to advantage for other young pulsars which have high timing noise.

(2) We determined an accurate position for PSR J0205+6449 of 02°55′37″92, 64°49′41″3, which is 1.5 arcsec different than the previously accepted one, which was based on an X-ray image. Re-examination of the X-ray data, however, reveals an error in the original reported position; the newly determined X-ray position reported here is consistent with the VLBI position. Furthermore, this position is coincident with an optical source identified by Shearer & Neustroev (2008).

(3) We determined the proper motion of PSR J0205+6449. After correction for Galactic rotation, we found a proper motion of \( 2.3 \pm 0.3 \) mas yr\(^{-1}\), corresponding to tangential velocity of \( 35 \pm 6 \) (\( D/3.2 \) kpc) km s\(^{-1}\) at PA \( \sim 38^\circ \). This low speed puts PSR J0205+6449 amongst the slowest few percent of young pulsars.

(4) We estimated PSR J0205+6449’s position at birth. If it is the remnant of a supernova in 1181 AD, then its position at birth is only \( 1.9 \) arcsec different than at present. If as seems more likely, the age of PSR J0205+6449 and 3C 58 is several thousand years, then its position at birth was near the midpoint of the presently visible nebula, although somewhat displaced from the present radio brightness centre which is close to the present location of PSR J0205+6449.

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