THE FLUX VARIABILITY OF MARKARIAN 501 IN VERY HIGH ENERGY GAMMA RAYS

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Received 1998 December 8; accepted 1999 January 22

ABSTRACT

The BL Lacertae object Markarian 501 was identified as a source of γ-ray emission at the Whipple Observatory in 1995 March. Here we present a flux variability analysis on several timescales of the 233 hr data set accumulated over 213 nights (from March 1995 to July 1998) with the Whipple Observatory 10 m atmospheric Cerenkov imaging telescope. In 1995, with the exception of a single night, the flux from Markarian 501 was constant on daily and monthly timescales and had an average flux of only 10% that of the Crab Nebula, making it the weakest very high energy source detected to date. In 1996, the average flux was approximately twice the 1995 flux and showed significant month-to-month variability. No significant day-scale variations were detected. The average γ-ray flux above ~350 GeV in the 1997 observing season rose to 1.4 times that of the Crab Nebula—14 times the 1995 discovery level—allowing a search for variability on timescales shorter than 1 day. Significant hour-scale variability was present in the 1997 data, with the shortest, observed on MJD 50,607, having a doubling time of ~2 hr. In 1998 the average emission level decreased considerably from that of 1997 (to ~20% of the Crab Nebula flux), but two significant flaring events were observed. Thus the emission from Markarian 501 shows large amplitude and rapid flux variability at very high energies, as does Markarian 421. It also shows large mean flux level variations on year-to-year timescales, behavior that has not been seen from Markarian 421 so far.

Subject headings: BL Lacertae objects: individual (Markarian 501) — gamma rays: observations

1. INTRODUCTION

Three active galactic nuclei (AGNs) have been discovered to be very high energy (VHE; E ∼ 300 GeV) γ-ray sources by the Whipple Observatory γ-ray collaboration: Markarian 421 (Punch et al. 1992), Markarian 501 (Quinn et al. 1996), and 1ES 2344+514 (Catanesi et al. 1998). These are the three closest BL Lacertae objects (BL Lacs), with redshifts in the range 0.0308–0.044, and are among the brightest at X-ray energies. They are all classified as X-ray–selected BL Lacs (XBLs), as their synchrotron spectra extend into the X-ray range. A fourth BL Lac, PKS 2155–304, has been detected in VHE γ-rays by the University of Durham group (Chadwick et al. 1999).

The Energetic Gamma-Ray Experiment Telescope (EGRET) on board the Compton Gamma Ray Observatory (CGRO) has detected at least 51 AGNs at energies greater than 100 MeV (Thompson et al. 1995; Mukherjee et al. 1997). They are all members of the blazar class of AGNs, which include flat spectrum radio quasars and BL Lacs. Of the EGRET-detected blazars, 14 are BL Lacs, with 12 radio-selected BL Lacs and only two X-ray–selected BL Lacs. Mrk 421 is the only VHE source in this catalog, and it is among the weakest. However, Mrk 501 has recently been detected at the 4σ level with EGRET (Kataoka et al. 1999). One of the most striking characteristics of the EGRET-detected blazars is variability; 42 of the 51 AGNs exhibit variability (Mukherjee et al. 1997). Variability timescales as short as 4 hr have been observed (Mattox et al. 1997).

VHE observations of Mrk 421 and Mrk 501 have also revealed extreme variability (see, e.g., Gaidos et al. 1996; Quinn et al. 1996). The VHE flux has been measured to vary by nearly a factor of 100 in Mrk 421 (McEnery et al. 1999) and, as we show in the following sections, Mrk 501 has been measured with fluxes ranging from 0.1 to 5 times the Crab Nebula flux with the Whipple Observatory atmospheric Cerenkov telescope. The large collection area (∼3.5 × 105 m2) of the Whipple Observatory telescope permits sensitive studies of variability on timescales inaccessible to space-based telescopes. Indeed, the shortest observed variability of any blazar at any γ-ray energy, a 30 minute duration flare observed from Mrk 421 (Gaidos et al. 1996), was measured with this telescope. Both Mrk 421 and Mrk 501 have been closely monitored by the Whipple Collaboration since their discovery, with an ∼0.5 hr exposure per night being sufficient for detection of flaring activity. Prior to 1997, the γ-ray emission from Mrk 421 was generally observed to have a higher mean flux (Schubnell et al. 1996) and to have been more frequently variable (Buckley et al. 1996) than that of Mrk 501.

In the spring of 1997, the Whipple Collaboration observed Mrk 501 to be in an unprecedented high-emission...
state at VHE energies, as subsequently confirmed by several independent Cerenkov imaging groups (Protheroe et al. 1998). A public CGRO target of opportunity was initiated in response to a request by the Whipple Collaboration in 1997. Evidence of correlated variability in data from the Whipple Observatory γ-ray telescope, the Oriented Scintillation Spectrometer Experiment (OSSE) and the All-Sky Monitor (ASM) of the Rossi X-Ray Timing Explorer is presented in Catanese et al. (1997). During this observing campaign, the energy output of Mrk 501 in VHE γ-rays was comparable to that in the 2–100 keV range, but the variability amplitude was larger. The correlations seen may imply some relativistic beaming of the emission, given that the variability amplitude was larger. The correlations seen may imply some relativistic beaming of the emission, given that the variability amplitude was larger.

During 1995 and 1996 the camera consisted of 2, and the analysis methodology, including γ-ray selection criteria and the test for variability, is described in § 3. The results of the analysis are presented in § 4, and their implications are briefly discussed in § 5.

2. SOURCE OBSERVATIONS

Observations were made with the Whipple Observatory 10 m atmospheric Cerenkov imaging telescope (Cawley et al. 1990), located on Mount Hopkins in southern Arizona. The 10 m reflector images the Cerenkov radiation from cosmic-ray- and γ-ray–initiated air showers onto a high-resolution camera mounted in the focal plane. Subsequent off-line analysis of the images (described below) facilitates the selection of candidate γ-ray events.

The high-resolution camera utilizes fast photomultiplier tubes (PMTs) arranged in a hexagonal array, with intertube spacing of 0.25. During 1995 and 1996 the camera consisted of 109 PMTs, with a resulting field of view (FOV) of ~2.8”. An event was recorded when any two of the inner 91 PMTs registered a signal ≥40 photoelectrons within an effective resolving time of 15 ns. For the 1997 observations a further 42 PMTs were added, resulting in an FOV of ~3.4’. The trigger condition remained the same as for the 109 PMT camera. In the summer of 1997, a new camera containing 331 pixels was installed. This camera has an FOV of ~4.8”. The trigger condition for this enlarged camera required that any two of the 331 pixels produce a signal ≥40 photoelectrons within an effective resolving time of 8 ns. The telescope was also triggered artificially, after every 24 events for the 109/151 pixel cameras and once every second for the 331 pixel camera, to determine the background sky-brightness level in each PMT. Light cones, which minimize the dead space between PMTs and reduce the albedo effect, were used on the 109 and 151 pixel cameras, but were not yet installed on the 331 pixel camera.

In general, two modes of observation are used: ON/OFF and TRACKING. With the ON/OFF mode, the source is tracked continuously for 28 minutes and then, to estimate the background, a region offset in right ascension (R.A.) by 30 minutes (allowing 2 minutes sleww time) is tracked. This has the disadvantage that an equivalent amount of observation time is spent looking away from the source. Alternatively, the TRACKING mode, where the background is estimated from the ON-source run itself (see § 3.2), can be used. In this case an OFF-source run is not required for each ON-source run, allowing continuous monitoring of an object. However, OFF runs are still needed to determine the response of the telescope to background events.

Observations are typically made when the source zenith angle is less than 35° and are referred to as small zenith angle (SZA) observations. Observations at large zenith angles (LZA, typically 55°–70°) may also be made. Increasing the zenith angle has the effect of increasing the energy threshold, but it has the benefit of increasing the collection area. Thus it is an excellent method for increasing photon statistics to facilitate the determination of the energy spectrum at higher energies. For a detailed description of the LZA technique, see Krennrich et al. (1997).

Since its discovery as a γ-ray source in 1995 (Quinn et al. 1996), the VHE γ-ray emission from Mrk 501 has been monitored intensively with the Whipple Observatory 10 m telescope. Only SZA observations taken under good sky conditions have been considered in the analysis for variability presented here. Our selection includes data from 56, 50, 55, and 49 nights of observation in the spring–summer periods of 1995, 1996, 1997, and 1998, respectively. Table 1 summarizes the resulting database.

3. DATA ANALYSIS

3.1. Gamma-Ray Selection

The vast majority of events detected by Cerenkov telescopes are cosmic rays. Candidate γ-ray events are selected on the basis of the shape and orientation of the Cerenkov images: γ-ray images are typically more compact and elliptical than background hadronic images and tend to have their major axes aligned with the source location in the FOV. Background cosmic-ray events, on the other hand, have random orientations.

Each image is first subjected to a cleaning procedure (Fegan 1997) that suppresses pixels that are dominated by light from fluctuations of the night-sky background. A moment-fitting routine is then used to calculate various image parameters. The shape of each image is characterized by the parameters length and width and the orientation by θ, the angle between the major axis of the image and the line joining the source location in the FOV to the centroid of the image. In addition, for the data taken in 1998, an asymmetry cut has been included. γ-ray images have a cometary shape with a tail that points away from the source location in the FOV (Buckley et al. 1998), and thus their intensity profiles have positive asymmetry. The larger FOV of the camera used in 1998 allows this parameter to be accurately determined, something which was not possible with the smaller FOV cameras.

| Month   | 1995 | 1996 | 1997 | 1998 |
|---------|------|------|------|------|
| Feb     |      |      | 3.3  | 0.9  |
| Mar     | 3.7  | 18.0 | 12.4 | 5.5  |
| Apr     | 17.7 | 15.5 | 20.7 | 4.6  |
| May     | 20.3 | 10.3 | 21.2 | 20.5 |
| Jun     | 11.5 | 8.8  | 9.7  | 11.6 |
| Jul     | 8.3  |      |      | 8.7  |
| Total   | 61.5 | 52.6 | 67.3 | 51.8 |

* Combined hours of on and tracking data.
Before the application of shape and orientation cuts, a software trigger cut is also applied to eliminate events close to threshold, some of which are induced by noise fluctuations. The software trigger involves cuts on the image size (i.e., the total number of photoelectrons recorded), the counts in each of the brightest two tubes (MAX1, MAX2), and a requirement that at least three tubes above a low noise threshold (2.25 $\sigma$, where $\sigma$ is the rms sky noise in a PMT, as determined from the artificially triggered events) be neighbors (NBR 3). A distance cut is applied to eliminate images that are too close to the center of the camera and will have poor $\alpha$ reconstruction, and also those events that have occurred too close to the edge of the FOV and may be truncated.

Due to the continuous evolution of the high-resolution camera, and changes such as deterioration of mirror reflectivity through weathering and the presence or absence of light cones, the optimum data analysis cuts differ for each year. The cuts used for a given telescope configuration are optimized on an independent data set, usually data taken on the Crab Nebula or, in the case of the 1998 data, Mrk 421. Table 2 lists the cuts used for each years’ analyses.

3.2. Gamma-Ray Rate and Flux Calculation

There are slightly different analysis methods for the on/off and tracking observation modes. For the on/off observations, the background is estimated from the off-source run, which is assumed to be on a sky region that does not include a gamma-ray source. This analysis mode has been discussed at length elsewhere (see, e.g., Kerrick et al. 1995; Catanese et al. 1998).

For tracking observations the background is estimated from the on-source run itself. All of the gamma-ray selection criteria apart from orientation ($\alpha$) are applied to the data. The background is then estimated from events that are not oriented toward the source. In this analysis, background events with values of $\alpha$ between 20° and 65° are used. Images having values of $\alpha$ between 65° and 90° are discarded because of possible systematic effects due to truncation at the camera’s edge. Once the number of events with orientations in the 20°–65° range is known, then the expected number of background events in the signal domain ($\alpha$ of 0°–15°, or to 10° for data taken in 1998) can be estimated. Off-source data recorded for this source and others can be combined to calculate a ratio, $R \pm \Delta R$, of the number of events in the signal region to those in the 20°–65° region in the absence of a source. In this case, the significance of a gamma-ray excess, $S$, is given by

$$S = \frac{N_{on} - rN_{off}}{\sqrt{N_{on} + r^2 N_{off} + (\Delta r)^2 N_{off}^2}},$$

(1)

and the gamma-ray rate ($R \pm \Delta R$) is calculated from

$$R \pm \Delta R = \frac{N_{on} - rN_{off}}{t} \pm \sqrt{\frac{N_{on} + r^2 N_{off} + (\Delta r)^2 N_{off}^2}{t}},$$

(2)

where $N_{on}$ is the number of counts in the gamma-ray domain ($\alpha < 10°$ or 15°), $N_{off}$ is the number of counts in the 20°–65° $\alpha$ range, and $t$ is the duration of the observation.

The inclusion of the statistical error on the tracking ratio effectively limits the amount (duration) of tracking data that can be usefully analyzed for an excess. Once the duration of the tracking data exceeds that of the off-source data used in the calculation of the tracking ratio, the error on the tracking ratio starts to dominate and to limit the significance of a detection.

For the purpose of investigating possible differences in results produced by the on/off and tracking analyses, the Crab Nebula data were analyzed with both methods (Quinn 1997). The results demonstrated that the gamma-ray rates derived using both methods were consistent and stable; the tracking analysis did not introduce any apparent variability in the rate. For the analysis presented in this paper, all of the on-source data were combined with the tracking data, and the resulting database was analyzed using the tracking analysis.

The data presented here were taken with different telescope configurations having different sensitivities. It is therefore necessary to normalize when comparing the different data sets. To a first approximation, this can be achieved by converting the rates to fractions of the gamma-ray rate from the Crab Nebula taken with the same telescope configuration and analyzed with the same cuts. The Crab Nebula is believed to be a steady source of VHE gamma-rays, as has been observed by the Whipple Collaboration over the past decade (Cavley et al. 1999). To convert a given gamma-ray rate to an integral flux, the rate as a fraction of the Crab Nebula was multiplied by (1.05 ± 0.24) $\times 10^{-10}$ cm$^{-2}$ s$^{-1}$, which represents the integral Crab Nebula flux above 350 GeV (Hillas et al. 1998), the threshold of the analysis presented here.

3.3. Test for Variability

To search for temporal variability in the gamma-ray flux, we apply a $\chi^2$ test for a constant rate. The $\chi^2$ sum is converted into a probability ($P_{\chi^2}$) that the emission is constant about the mean using the incomplete gamma function $\text{gammi}(a, x)$.
(Press et al. 1988),

\[ P_{\chi^2} = \text{gamma}\left(\frac{N-1}{2}, \frac{\chi^2}{2}\right), \]

where \( N - 1 \) is the number of degrees of freedom. The number of trials is taken into account by calculating the probability \( P_{\chi^2} \) of \( \chi^2 \) occurring in \( N \) trials from

\[ P_{\text{trials}} = 1 - (1 - P_{\chi^2})^N. \]

This method is used to test whether the distribution of measured \( \gamma \)-ray rates is consistent with statistical fluctuations about the mean for a range of timescales.

4. RESULTS

The data have been analyzed using the TRACKING analysis described above. This differs from that applied to the 1995 data by Quinn et al. (1996), in that a statistical error on the TRACKING ratio is now included and that a 10% systematic error is no longer added to TRACKING results, as a careful study showed that the results of the ON/OFF and TRACKING analysis methods are in close agreement (Quinn 1997). In fact, these changes tend to cancel each other.

For the 1995 data set as a whole, we obtain a \( \gamma \)-ray rate of approximately 10% of that of the Crab Nebula \((0.18 \pm 0.02 \text{ minute}^{-1})\), giving a 9.1 \( \sigma \) excess. This rose to approximately 20% of the Crab Nebula rate for the following season (analysis of 1996 data reveals an excess of 11.1 \( \sigma \) and a \( \gamma \)-ray rate of \(0.26 \pm 0.02 \text{ minute}^{-1}\)), indicating that the average emission level had doubled since the previous year.

The monthly and nightly average rates, in fractions of the Crab Nebula rate, over all 4 yr of Mrk 501 observation are shown in Figure 1. The rate in 1995 appears to have been constant with the exception of one night, MJD 49,920, when the emission level had doubled since the previous year.

The emission appeared to increase steadily from 1997 February through May and then to level off in 1997 June. The daily rates also exhibit dramatic variability. Day-to-day changes in the flux by factors greater than 4 were observed, and on eight occasions the flux more than doubled between consecutive nights. On four occasions there were equally rapid decays in the rate. The average flux level for the season is 1.4 times that of the Crab Nebula—an increase by a factor of 14 from the level in 1995. The peak rate, observed on MJD 50,554, is 3.7 times the Crab rate. The average rates for the 1997 data were calculated from only one run (the first with elevation above 55\(^\circ\)) per night. This was done to remove a bias in the observing strategy whereby observations of Mrk 501 on a given night continued only if the source was in a very active state.

The strong signal-to-background level in the 1997 data allowed a search to be made for variability on timescales shorter than 1 day. For this test, data from each night on which there were three or more runs (approximately 1.5 hr, see § 2) were analyzed to test for run-to-run variability. A total of 24 nights satisfied this criterion. The \( \chi^2 \) probability that the emission was constant was calculated for each of

| Year | Number of Months | Rate\(^a\) | \( P_{\text{trials}} \)\(^b\) |
|------|-----------------|----------|-----------------|
| 1995 | 5               | 0.10 ± 0.01 | 0.05 |
| 1996 | 4               | 0.20 ± 0.02 | 3.6 × 10\(^{-4}\) |
| 1997 | 5               | 1.39 ± 0.07 | 8.8 × 10\(^{-2}\) |
| 1998 | 6               | 0.32 ± 0.04 | 8.5 × 10\(^{-1}\) |

\(^a\) Average rate over the observation period expressed as a fraction of the rate from the Crab Nebula.

\(^b\) \( \chi^2 \) probability for constant emission after accounting for trials.
The nights; the distribution of these probabilities is shown in Figure 2. For a statistically variable source this distribution should be flat, but there is an excess of nights with small probabilities. Assuming that the variations are purely statistical, the probability of getting seven nights in this first bin (width equal to 0.025) out of 24 trials is $1.45 \times 10^{-3}$. Thus identify these two nights as exhibiting significant hour-scale variability. The flux on MJD 50,607 has a doubling time of ~10^-6, while for MJD 50,607 the probability is $5.8 \times 10^{-8}$ (after accounting for trials). The flux on MJD 50,607 has a doubling time of ~2 hr. We thus identify these two nights as exhibiting significant hour-scale variability, while the five other nights exhibit marginal variability. Figure 3 shows the γ-ray rates for the two nights with significant variability.

A search of the 1997 data for variability on timescales of less than half an hour has also been performed. For this test each of the 28 minute runs, 144 in total, was divided into three equal length intervals. Each triplet was then analyzed for variability. No significant variations were found and the distribution of $\chi^2$ probabilities (Fig. 2b) does not indicate any excess of low probabilities. Hence, we see no evidence for significant subhour-scale variability. The average γ-ray rate for the 1998 data set was $0.42 \pm 0.04$ minute^{-1}, approximately 20% of the rate obtained from the Crab Nebula, i.e., on average the emission was much lower than in 1997. In fact, the average rates for 1998 March, April, and May are comparable to the level of the initial detection in 1995. There were, however, two significant flaring events. The first occurred in early 1998.

### Table 4
RESULTS OF ANALYSIS FOR VARIABILITY OF THE DAILY AVERAGES

| Observation Period   | Number of Nights | Rate | $P_{\text{trials}}$  |
|----------------------|------------------|------|----------------------|
| 1995                 |                  |      |                      |
| Mar 28–Apr 07        | 09               | 0.09±0.02 | 1.00                |
| Apr 22–May 10        | 15               | 0.08±0.02 | 1.00                |
| May 21–Jun 05        | 14               | 0.07±0.02 | 1.00                |
| Jun 18–Jul 01        | 10               | 0.10±0.03 | 0.99                |
| Jul 22–Jul 30        | 08               | 0.16±0.02 | $1.8 \times 10^{-3}$ |
| Total                | 56               | 0.10±0.01 | 0.22                |
| 1996                 |                  |      |                      |
| Mar 17–Mar 30        | 13               | 0.19±0.02 | 0.33                |
| Apr 13–Apr 28        | 15               | 0.12±0.02 | 1.00                |
| May 09–May 25        | 15               | 0.24±0.03 | 0.18                |
| Jun 06–Jun 21        | 10               | 0.29±0.03 | 0.46                |
| Total                | 53               | 0.26±0.02 | $3.8 \times 10^{-6}$ |
| 1997                 |                  |      |                      |
| Feb 10–Feb 16        | 05               | 0.62±0.08 | 0.33                |
| Mar 05–Mar 16        | 12               | 0.96±0.08 | $<10^{-38}$         |
| Apr 07–Apr 19        | 12               | 1.37±0.08 | $<10^{-38}$         |
| Apr 30–May 15        | 12               | 1.68±0.09 | $<10^{-38}$         |
| May 27–Jun 11        | 14               | 1.61±0.09 | $<10^{-38}$         |
| Total                | 55               | 1.29±0.06 | $<10^{-38}$         |
| 1998                 |                  |      |                      |
| Feb 28–Mar 09        | 06               | 1.94±0.29 | $<10^{-38}$         |
| Mar 23–Mar 31        | 03               | 0.00±0.07 | 1.00                |
| Apr 19–Apr 29        | 09               | 0.08±0.04 | 1.00                |
| May 20–May 31        | 10               | 0.09±0.02 | 0.13                |
| Jun 16–Jul 03        | 15               | 0.30±0.05 | $<10^{-38}$         |
| Jul 15–Jul 27        | 06               | 0.12±0.04 | 0.94                |
| Total                | 49               | 0.32±0.04 | $<10^{-38}$         |

*a* Average rate over the observation period expressed as a fraction of the rate from the Crab Nebula.

*b* $\chi^2$ probability for constant emission after accounting for trials.
March where an apparent rise and decay were observed. Unfortunately, this flare is poorly sampled due to bad weather. The flux was observed to be relatively high (~1.3 times the Crab Nebula flux) on MJD 50,876, and on the following night a flux of ~5.0 times that of the Crab Nebula was recorded. This is the largest flux detected to date from Mrk 501 by the Whipple Observatory 10 m telescope. For the next observation on MJD 50,880, the measured flux was still relatively high (~1.3 times the level of the Crab Nebula flux). A second flare occurred in 1998 June. The flux increased on two consecutive nights, peaking at ~1.1 times the Crab Nebula flux on MJD 50,991, then decayed on a similar timescale. The average flux level immediately to either side of this flare was below the sensitivity of the telescope.

5. DISCUSSION

We have demonstrated that rapid variability, a common characteristic of blazars at all observed energies, is also present in the VHE γ-ray emission from Mrk 501. Our 4 yr data set spans a remarkable change in the flux level from Mrk 501: the average yearly emission level exhibited a 14-fold increase between 1995 and 1997, and the average daily flux varied by a factor of ~50 (see Fig. 1). In 1997 large amplitude day-scale flares occurred frequently and were usually followed by equally rapid decays. Day-scale changes in flux by factors as large as 4.7 were observed. Episodes of significant hour-scale variability were detected, with one having a doubling time of 2 hr. In addition, there is evidence of consistent hour-scale variability which is not resolved in individual episodes. A variable flux from Mrk 501 was reported by at least four other atmospheric Cerenkov observatories in 1997 (see, e.g., Protheroe et al. 1998).

The data presented here suggest that Mrk 501 was more variable when the flux level was higher. However, this effect could also be due to the sensitivity of the telescope. At low flux levels it takes longer to accumulate a significant signal, so the search for short-term variability may be limited by poor statistics. To address this issue we performed a test to see if the day-scale variability observed in 1997 would have been detected in 1996 and/or 1995, and if the month-scale variability seen in 1996 and 1997 would have been detected in 1995. We calculated the percentage deviations about the mean level from another period, calculated the signal (and statistical error) that would have been observed given the same percentage deviations about that mean. The results are that neither the 1997 month-scale nor day-scale variability would have been detected at a significant level if present in 1995. However, the 1996 month-scale variability would have been significant if present in 1995 (chance probability of ~10⁻⁴), while the 1997 degree of day-scale variability, if present in 1996, would have been detectable (chance probability ~10⁻⁷). Thus we conclude that there was a change in the flaring characteristics, in addition to the change in the mean flux level between the different observing seasons.

The VHE γ-ray emission from Mrk 501 exhibits rapid variability similar to that of the emission from Mrk 421. A major difference between the two objects is that the emission from Mrk 501 seems to have a base level, which changes on monthly and yearly timescales, whereas the VHE γ-ray emission from Mrk 421 has been described as consisting of a series of rapid flares with no underlying baseline (Buckley et al. 1996). The variability of the VHE γ-ray emission of Mrk 501 in 1995 and 1996 was similar to that at other wavelengths, with small amplitude, slow variations being more common than fast, large amplitude flares.

The increase in the VHE γ-ray power and variability in 1997 was accompanied by an increase in the hard X-ray power and an extension of the synchrotron spectrum to at least 100 keV (Catanese et al. 1997). This is consistent with an inverse Compton mechanism for producing VHE γ-rays. However, our observations of Mrk 501 to date cannot discriminate between the inverse-Compton (electron) models and those where the dominant particles producing γ-rays in the jet are protons. More densely sampled light curves, covering as broad a wave band as possible, are needed to provide more insight into the mechanisms responsible for the VHE radiation from BL Lac objects. Future, more sensitive γ-ray observations of AGNs with proposed detectors such as GLAST, HESS, MAGIC, and VERITAS will allow the structure of flares on shorter timescales to be determined.

We acknowledge the technical assistance of K. Harris and E. Roache. This research is supported by grants from the US Department of Energy and by NASA, PPARC in the UK, and Forbairt in Ireland.

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