On Barnes Beta Distributions and Applications to the Maximum Distribution of the 2D Gaussian Free Field

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Abstract

A new family of Barnes beta distributions on $(0, \infty)$ is introduced and its infinite divisibility, moment determinacy, scaling, and factorization properties are established. The Morris integral probability distribution is constructed from Barnes beta distributions of types $(1,0)$ and $(2,2)$, and its moment determinacy and involution invariance properties are established. For application, the maximum distributions of the 2D gaussian free field on the unit interval and circle with a non-random logarithmic potential are conjecturally related to the critical Selberg and Morris integral probability distributions, respectively, and expressed in terms of sums of Barnes beta distributions of types $(1,0)$ and $(2,2)$.

Keywords: Infinite divisibility, multiplicative chaos measure, Morris integral, multiple gamma function, Selberg integral, self-duality.

1 Introduction

This paper advances the theory of Barnes beta probability distributions that we introduced in [38] and [39]. Barnes beta distributions lie at the crossroads of several areas of probability theory and statistical physics: Dufresne distributions [11], infinite divisibility in the context of special functions of analytic number theory [5], [23], [35], Lévy processes [26] – [28], and, conjecturally, gaussian multiplicative chaos [37], the distribution of the maximum of the characteristic polynomial of GUE random matrices [21], and mod-gaussian limit theorems for mesoscopic statistics [40].

The defining property of Barnes beta distributions $\beta_{M,N}(a,b)$, $M \leq N$, $a = (a_1, \cdots, a_M)$, $b = (b_0, b_1, \cdots, b_N)$, is that their Mellin transform is given in the form of an intertwining product of ratios of multiple gamma functions of Barnes [4]. For example, in the special case of type $(2,2)$, $a = (a_1, a_2)$, $b = (b_0, b_1, b_2)$, the Mellin transform is

$$\mathbb{E}[\beta_{2,2}(a,b)^q] = \frac{\Gamma_2(q|a)}{\Gamma_2(b_0|a)} \frac{\Gamma_2(b_0+b_1|a)}{\Gamma_2(q+b_0+b_1|a)} \frac{\Gamma_2(b_0+b_2|a)}{\Gamma_2(q+b_0+b_2|a)} \times \frac{\Gamma_2(q+b_0+b_1+b_2|a)}{\Gamma_2(b_0+b_1+b_2|a)}.$$

It is known that $\beta_{M,N}(a,b)$ takes values in $(0, 1]$ if $M < N$ and in $(0, 1)$ if $M = N$.

The contribution of this paper is to review the general theory of Barnes beta distributions for $M \leq N$, extend it to the case of $M = N + 1$, and give novel applications of the new and existing...
theory. Our theoretical contribution is the proof that \( \beta_{M,M-1}(a,b) \) exists for \( M \in \mathbb{N} \), takes values in \((0,\infty)\), has infinitely divisible logarithm, and satisfies remarkable infinite factorizations. In particular, the new construction specializes to certain intertwining products of multiple sine functions. This is motivated by the work of \([26]-[28]\), who computed the Mellin transform of certain functionals of the stable Lévy process in the form of products of ratios of double gamma functions that are different from \(1\). Their distributions coincide with what we call \( \beta_{2,1}(a,b), a = (a_1,a_2), b = (b_0,b_1), \)

\[
E[\beta_{2,1}(a,b)^q] = \frac{\Gamma_2(q+b_0|a)}{\Gamma_2(b_0|a)} \frac{\Gamma_2(b_0+b_1|a)}{\Gamma_2(q+b_0+b_1|a)}.
\]

The contribution of this paper to statistical physics is to formulate precise conjectures about the distribution of the maximum of the 2D gaussian free field (GFF) on the unit interval and circle with a non-random logarithmic potential in terms of Barnes beta distributions. The GFF is a fascinating mathematical object that is of fundamental interest in statistical mechanics of disordered energy landscapes \([7,14,16,18,19,20]\), quantum disordered systems exhibiting multifractality \([8,9,15]\), and quantum gravity \([12,42]\). It also appears naturally in the limit of a wide class of statistics. For example, \([22]\) proved convergence of the distribution of the log-characteristic polynomial of large unitary matrices to the GFF on the circle and \([17]\) proved convergence of an appropriately rescaled log-characteristic polynomial of large Hermitian matrices to the GFF on the interval. We recently showed in \([40]\) that the smoothed indicator function of the mesoscopic statistics of Riemann zeroes of Bourgade-Kuan \([6]\) and Rodgers \([43]\) converges to the GFF on the interval. We also indicated there that the same result is true of any linear statistic that converges to a gaussian process having \(\mathcal{H}^{1/2}(\mathbb{R})\) limiting covariance such as the local CUE statistic of Soshnikov \([46]\), for example.

The problem of calculating the maximum distribution of the GFF was first considered in \([14]\) for the GFF on the circle and in \([19]\) on the interval, and then extended in \([15,20]\), and most recently in \([7,18]\). In all cases they conjectured the Laplace transform of the maximum distribution including the leading asymptotic (non-random) drift and the fluctuating part of the distribution. While there has recently been made good progress in verifying their results for the drift, see \([10]\), their conjectures for the fluctuating part are still beyond the reach of existing methods, as the law of the so-called derivative martingale is unknown, see \([30,48]\) for details. The difficulty of the problem is that their solutions are based on the still unproven freezing scenario for the gaussian multiplicative chaos measure, see \([31]\) for recent progress, and on the conjectured form of the Mellin transform of the total (random) mass of that measure, see \([49]\) for partial results in the circle case. Our contribution is to show that the result of \([18,19]\) for the fluctuating part on the interval does correspond to a valid probability distribution, namely, the critical Selberg integral probability distribution, which we developed in \([37,39]\), and can be naturally expressed in terms of Barnes beta distributions of types \((1,0)\) and \((2,2)\), thereby conjecturing the law of the derivative martingale. We also clarify the origin of the self-duality of the Mellin transform that plays a key role in their freezing framework and give a heuristic derivation of the conjecture by combining their calculations for the fluctuating part with our conjecture about the total mass of the multiplicative chaos measure on the interval. In the case of the circle, we construct a new probability distribution having the property that its \(n\)th moment equals the value of the Morris integral of dimension \(n\), express it in terms of Barnes beta distributions of types \((1,0)\) and \((2,2)\), prove self-duality of its Mellin transform, and further extend the original work of \([14]\) by formulating a novel conjecture about (the fluctuating part of) the distribution of the maximum of the GFF on the circle with a non-random logarithmic potential in terms of the critical Morris integral distribution, thereby also conjecturing the law of the derivative martingale in this case.
The plan of the paper is as follows. In Section 2 we give a review of the general theory of Barnes beta distributions of type \((M,N)\) and of the Selberg integral distribution as a special case of the theory corresponding to type \((2,2)\). In Section 3 we construct a new family of Barnes beta distributions on \((0,\infty)\). In Section 4 we construct the Morris integral distribution and study its properties. In Section 5 we treat the maximum distribution of the gaussian free field. In Section 6 we give the proofs. Section 7 concludes.

2 A Review of Multiple Gamma Functions and Barnes Beta and Selberg Integral Distributions

The multiple gamma function of Barnes [4] is defined classically by

\[
\Gamma^{-1}_M(w|a) = \mathcal{P}(w|a) w \prod_{n_1,\ldots,n_M=0}^{\infty} \left(1 + \frac{w}{\Omega}\right) \exp\left(\sum_{k=1}^{M} \frac{(-1)^k}{k} \frac{w^k}{\Omega^k}\right),
\]

where \(\Re(w) > 0\), \(\mathcal{P}(w|a)\) is a polynomial in \(w\) of degree \(M\) that depends on one’s choice of normalization, the parameters \(a_j > 0\), \(j = 1 \cdots M\),

\[
\Omega = \sum_{i=1}^{M} n_i a_i,
\]

and the prime indicates that the product is over all indices except \(n_1 = \cdots = n_M = 0\). Our choice of normalization is that of Ruijsenaars [44], which is explained next. Let

\[
f_M(t|a) = t^M \prod_{j=1}^{M} \left(1 - e^{-at}\right)^{-1},
\]

for some integer \(M \geq 0\). Slightly modifying the definition in [44], we define multiple Bernoulli polynomials by

\[
B_{M,m}(x|a) \triangleq \frac{d^m}{dt^m} \bigg|_{t=0} \left[ f_M(t|a) e^{-xt} \right].
\]

The key result of [44] about the log-multiple Gamma function that we need is summarized in the following theorem.

**Theorem 2.1 (Ruijsenaars)** \(\log \Gamma_M(w|a)\) satisfies the Malmstén-type formula for \(\Re(w) > 0\),

\[
\log \Gamma_M(w|a) = \int_0^\infty \frac{dt}{t^{M+1}} \left( e^{-wt} f_M(t|a) - \sum_{k=0}^{M-1} \frac{t^k}{k!} B_{M,k}(w|a) - \frac{t^M}{M!} B_{M,M}(w|a) \right),
\]

\[
\log \Gamma_M(w|a) \text{ satisfies the asymptotic expansion,}
\]

\[
\log \Gamma_M(w|a) = -\frac{1}{M!} B_{M,M}(w|a) \log(w) + \sum_{k=0}^{M} \frac{B_{M,k}(0|a)(-w)^{M-k}}{k!(M-k)!} \sum_{l=1}^{M-k} \frac{1}{l} + R_M(w|a),
\]

\[
R_M(w|a) = O(w^{-1}), |w| \to \infty, |\arg(w)| < \pi.
\]
The formula in (7) can be thought of as defining the Barnes multiple gamma function by
\[ \Gamma_M(w \mid a) \triangleq \exp \left( \log \Gamma_M(w \mid a) \right), \] (10)
in which case one proves that the formula in (3) holds for a particular choice of \( P(w \mid a) \). One can show that \( \Gamma_M(w \mid a) \) satisfies the fundamental functional equation
\[ \Gamma_M(w \mid a) = \Gamma_M(w + a_j \mid a) \Gamma_M(w + a_i \mid a), \quad i = 1 \cdots M, \quad M \in \mathbb{N}, \] (11)
where \( a = (a_1, \cdots, a_i, a_{i+1}, \cdots, a_M) \). For example,
\[ \Gamma_1(w \mid a) = \frac{e^{w(a - 1)/2}}{\sqrt{2\pi}} \Gamma(w), \] (12)
\[ \Gamma_0(w) = 1/w. \] (13)

The multiple gamma function has three additional properties that are important for our purposes.

Let \( \text{Re}(w) > 0, \kappa > 0 \) and \( (\kappa a)_i \triangleq \kappa a_i, \quad i = 1 \cdots M \). Then, it has the scaling property,
\[ \Gamma_M(\kappa w \mid \kappa a) = \kappa^{-B_M(w \mid a)/M!} \Gamma_M(w \mid a). \] (14)

Let \( \text{Re}(w) > 0 \) and \( k = 1, 2, 3, \cdots \). It has the multiplication property,
\[ \Gamma_M(kw \mid a) = k^{-B_M(kw \mid a)/M!} \prod_{j=1}^{M!} \Gamma_M(w + \sum_{j=1}^{M!} \kappa \mid a). \] (15)

Given \( x > 0 \) and \( a = (a_1 \cdots a_{M-1}) \), there exist functions \( \Phi_M(w, x \mid a, a_M) \) and \( \Psi_M(w, y \mid a) \) such that
\[ \Gamma_M(w \mid a, a_M) = e^{\Phi_M(w, x \mid a, a_M)} \Gamma_M(w \mid a) \prod_{k=1}^{\infty} \frac{\Gamma_M(w + kaM \mid a)}{\Gamma_M(w + x + kaM \mid a)} \times \exp \left( \Psi_M(x, kaM \mid a) - \Psi_M(w, kaM \mid a) \right). \] (16)

\( \Psi_M(w, y \mid a) \) and \( \Phi_M(w, x \mid a, a_M) \) are polynomials in \( w \) of degree \( M \). This is known as Shintani factorization, see [45] for the original result for \( M = 2 \). The interested reader can find explicit formulas for \( \Psi_M(w, y \mid a) \) and \( \Phi_M(w, x \mid a, a_M) \) and derivations of all the three properties in [39].

We now proceed to review the Barnes beta construction following [38] and [39]. Define the action of the combinatorial operator \( \mathcal{S}_N \) on a function \( h(x) \) by

Definition
\[ (\mathcal{S}_N h)(q \mid b) \triangleq \sum_{p=0}^{N} (-1)^p \sum_{k_1 < \cdots < k_p = 1} h(q + b_{k_1} + \cdots + b_{k_p}). \] (17)

In other words, in (17) the action of \( \mathcal{S}_N \) is defined as an alternating sum over all combinations of \( p \) elements for every \( p = 0 \cdots N \).

Definition Given \( q \in \mathbb{C} - (-\infty, -b_0], \quad a = (a_1 \cdots a_M), \quad b = (b_0, b_1 \cdots b_N) \), let
\[ \eta_{M,N}(q \mid a, b) \triangleq \exp \left( (\mathcal{S}_N \log \Gamma_M)(q \mid a, b) - (\mathcal{S}_N \log \Gamma_M)(0 \mid a, b) \right). \] (18)

1 We will abbreviate \((\mathcal{S}_N \log \Gamma_M)(q \mid a, b)\) to mean the action of \( \mathcal{S}_N \) on \( \log \Gamma_M(x \mid a) \), i.e. \((\mathcal{S}_N \log \Gamma_M(x \mid a))(q \mid b)\).
The function $\eta_{M,N}(q|a, b)$ is holomorphic over $q \in \mathbb{C} - (-\infty, -b_0]$ and equals a product of ratios of multiple gamma functions by construction. Specifically,

$$\eta_{M,N}(q|a, b) = \frac{\Gamma_M(q + b_0|a)}{\Gamma_M(b_0|a) \prod_{j=1}^N \Gamma_M(b_0 + b_j|a)} \prod_{j=1}^N \frac{\Gamma_M(q + b_0 + b_j + b_{j_1}|a)}{\Gamma_M(b_0 + b_j + b_{j_1}|a)} \prod_{j_1 < j_2} \frac{\Gamma_M(b_0 + b_{j_1} + b_{j_2} + b_{j_1}|a)}{\Gamma_M(b_0 + b_{j_1} + b_{j_2}|a)} \cdots,$$

(19)

until all the $N$ indices are exhausted. The function $\log \eta_{M,N}(q|a, b)$ has an important integral representation that follows from that of $\log \Gamma_M(w|a)$ in (7).

**Theorem 2.2 (Existence and Structure)** Given $M, N \in \mathbb{N}$ such that $M \leq N$, the function $\eta_{M,N}(q|a, b)$ is the Mellin transform of a probability distribution on $(0, 1]$. Denote it by $\beta_{M,N}(a,b)$. Then,

$$\mathbb{E}[\beta_{M,N}(a,b)^q] = \eta_{M,N}(q|a, b), \quad \text{Re}(q) > -b_0. \quad (20)$$

The distribution $-\log \beta_{M,N}(a,b)$ is infinitely divisible on $[0, \infty)$ and has the Lévy-Khinchine decomposition for $\text{Re}(q) > -b_0$,

$$\mathbb{E}[\exp(q \log \beta_{M,N}(a,b))] = \exp\left(\int_0^\infty \left(1 - e^{-tq} - \frac{\prod_{j=1}^N (1 - e^{-b_0 t})}{\prod_{i=1}^M (1 - e^{-a_0 t})} dt \right)\right). \quad (21)$$

$\log \beta_{M,N}(a,b)$ is absolutely continuous if and only if $M = N$. If $M < N$,

$- \log \beta_{M,N}(a,b)$ is compound Poisson and

$$\mathbb{P}[\beta_{M,N}(a,b) = 1] = \exp\left(-\int_0^\infty e^{-b_0 t} \prod_{j=1}^N \frac{(1 - e^{-a_0 t})}{M} dt \right), \quad (22a)$$

$$= \exp\left(-\left(\mathcal{N}_M \log \Gamma_M\right)(0|a,b)\right). \quad (22b)$$

It is worth emphasizing that the integral representation of $\log \eta_{M,N}(q|a, b)$ in (21) is the main result as it automatically implies that $\beta_{M,N}(a,b)$ is a valid probability distribution having infinitely divisible logarithm, see Chapter 3 of [47] for background material on infinitely divisible distributions on $[0, \infty)$.

The Mellin transform of Barnes beta distributions satisfies a function equation that is inherited from that of the multiple gamma function and two remarkable factorizations.

**Theorem 2.3 (Properties)** $1 \leq M \leq N, q \in \mathbb{C} - (-\infty, -b_0], i = 1 \cdots M$,

$$\eta_{M,N}(q + a_i|a, b) = \eta_{M,N}(q|a, b) \exp\left(-\left(\mathcal{N}_M \log \Gamma_{M-1}\right)(q|\tilde{a}_i,b)\right). \quad (23)$$
Let \( \Omega \triangleq \sum_i n_i a_i \).

\[
\eta_{M,N}(q \mid a, b) = \prod_{k=0}^\infty \frac{\eta_{M-1,N}(q+ka_i \mid \hat{a}_i, b)}{\eta_{M-1,N}(ka_i \mid \hat{a}_i, b)}.
\]  

(24)

\[
\eta_{M,N}(q \mid a, b) = \prod_{n_1, \ldots, n_M=0}^\infty \left[ \frac{b_0 + \Omega}{q + b_0 + \Omega} \prod_{j_1=1}^N \frac{q + b_0 + b_{j_1} + \Omega}{b_0 + b_{j_1} + \Omega} \times \right.
\]

\[
\times \prod_{j_1 < j_2}^N \frac{q + b_0 + b_{j_1} + b_{j_2} + \Omega}{b_0 + b_{j_1} + b_{j_2} + \Omega} \times \left. \prod_{j_1 < j_2 < j_3}^N \frac{q + b_0 + b_{j_1} + b_{j_2} + b_{j_3} + \Omega}{b_0 + b_{j_1} + b_{j_2} + b_{j_3} + \Omega} \cdots \right].
\]  

(25)

Probabilistically, these factorizations are equivalent to, respectively,

\[
\beta_{M,N}(a, b) \overset{\text{inlaw}}{=} \prod_{k=0}^\infty \beta_{M-1,N}(\hat{a}_i, b_0 + ka_i, b_1, \ldots, b_N),
\]  

(26)

\[
\beta_{M,N}(a, b) \overset{\text{inlaw}}{=} \prod_{n_1, \ldots, n_M=0}^\infty \beta_{0,N}(b_0 + \Omega, b_1, \ldots, b_N).
\]  

(27)

We note that the factorizations in (24) and (25) correspond to the Shintani and Barnes factorizations of the multiple gamma function, see (16) and (3), respectively. The functional equation in (23) gives us the moments.

**Corollary 2.4 (Moments)** Assume \( a_i = 1 \). Let \( k \in \mathbb{N} \).

\[
\mathbb{E} \left[ \beta_{M,N}(a, b)^k \right] = \exp \left( -\sum_{l=0}^{k-1} \left( \mathcal{S}_N \log \Gamma_{M-1} \right)(l \mid \hat{a}_i, b) \right),
\]  

(28)

\[
\mathbb{E} \left[ \beta_{M,N}(a, b)^{-k} \right] = \exp \left( \sum_{l=0}^{k-1} \left( \mathcal{S}_N \log \Gamma_{M-1} \right)(-(l+1) \mid \hat{a}_i, b) \right), \quad k < b_0.
\]  

(29)

The scaling property in (14) gives us the scaling invariance.

**Theorem 2.5 (Scaling invariance)** Let \( \kappa > 0 \). Then,

\[
\beta_{M,N}^\kappa(\kappa a, \kappa b) \overset{\text{inlaw}}{=} \beta_{M,N}(a, b).
\]  

(30)

The interested reader can find additional properties of Barnes beta distributions and many examples in [38].

The multiplication property in (15) gives us the structure of the Selberg integral probability distribution, as was first shown in [39]. Recall the classical Selberg integral.

\[
\int \prod_{i=1}^l s_i^{\lambda_1} (1 - s_i)^{\lambda_2} \prod_{i<j} |s_i - s_j|^{-2/\tau} ds_1 \cdots ds_l =
\]

\[
\frac{\prod_{k=0}^{l-1} \Gamma(1 - (k + 1)/\tau) \Gamma(1 + \lambda_1 - k/\tau) \Gamma(1 + \lambda_2 - k/\tau)}{\Gamma(1 - 1/\tau) \Gamma(2 + \lambda_1 + \lambda_2 - (l + k - 1)/\tau)},
\]  

(31)
Theorem 2.6 (Selberg Integral Probability Distribution) Define the function

$$M(q | \tau, \lambda_1, \lambda_2) \triangleq \left( \frac{2\pi \tau^\tau}{\Gamma(1-1/\tau)} \right)^q \frac{\Gamma_2(1-q+\tau(1+\lambda_1)/\tau)}{\Gamma_2(1+\tau(1+\lambda_1)/\tau)} \times \frac{\Gamma_2(1-q+\tau(1+\lambda_2)/\tau)}{\Gamma_2(1+\tau(1+\lambda_2)/\tau)} \times \frac{\Gamma_2(2-q+\tau(2+\lambda_1+\lambda_2)/\tau)}{\Gamma_2(2-2q+\tau(2+\lambda_1+\lambda_2)/\tau)}$$

for $\text{Re}(q) < \tau$. Then, $M(q | \tau, \lambda_1, \lambda_2)$ is the Mellin transform of a probability distribution $M_{(\tau, \lambda_1, \lambda_2)}$ on $(0, \infty)$,

$$M(q | \tau, \lambda_1, \lambda_2) = \mathbb{E}[M_1^{(\tau, \lambda_1, \lambda_2)}]$$

and its positive moments satisfy for $l < \tau$

$$\mathbb{E}[M_1^{l(\tau, \lambda_1, \lambda_2)}] = \prod_{k=0}^{l-1} \frac{\Gamma(1-(k+1)/\tau)}{\Gamma(1-1/\tau)} \frac{\Gamma(1+\lambda_1-k/\tau)}{\Gamma(2+\lambda_1+\lambda_2-(l+k-1)/\tau)}.$$ 

Log $M_{(\tau, \lambda_1, \lambda_2)}$ is absolutely continuous and infinitely divisible. Define the distributions

$$L \triangleq \exp(A'(0, 4 \log 2/\tau)),$$ 

$$Y \triangleq \tau y^{-1} \exp(-y^{-\tau}) dy, y > 0,$$

i.e. log $L$ is a zero-mean normal with variance $4 \log 2/\tau$ and $Y$ is a power of the exponential. Let $X_1, X_2, X_3$ have the $\beta_{2,2}^{-1}(\tau, b)$ distribution with the parameters

$$X_1 \triangleq \beta_{2,2}^{-1}(\tau, b_0 = 1 + \tau, b_1 = \tau(\lambda_2 - \lambda_1)/2, b_2 = \tau(\lambda_2 - \lambda_1)/2),$$

$$X_2 \triangleq \beta_{2,2}^{-1}(\tau, b_0 = 1 + \tau, b_1 = \tau(\lambda_1 + \lambda_2)/2, b_2 = \tau(\lambda_2 - \lambda_1)/2),$$

$$X_3 \triangleq \beta_{2,2}^{-1}(\tau, b_0 = 1 + \tau, b_1 = 1 + \tau \lambda_1 + \tau \lambda_2/2, b_2 = 1 + \tau \lambda_1 + \tau \lambda_2/2).$$

Then,

$$M_{(\tau, \lambda_1, \lambda_2)} \xrightarrow{\text{inlaw}} 2\pi 2^{-\left[3(1+\tau) + 2\tau(\lambda_1 + \lambda_2)\right]/\tau} \frac{\Gamma(1-1/\tau)^{-1} L X_1 X_2 X_3 Y.}$$

The Mellin transform is involution invariant under

$$\tau \rightarrow \frac{1}{\tau}, q \rightarrow \frac{q}{\tau}, \lambda_i \rightarrow \tau \lambda_i.$$ 

$$\mathcal{M} \left( \frac{q}{\tau}, \frac{1}{\tau}, \tau \lambda_1, \tau \lambda_2 \right) (2\pi)^{-\frac{q}{\tau}} \Gamma^q(1-\tau) \Gamma(1-\tau)^{-q} \times \Gamma^q \left( \frac{1}{\tau} \right) \Gamma(1-q).$$
Remark The special case of $\lambda_1 = \lambda_2 = 0$ of Theorem 2.6 first appeared in [36]. The general case was first considered in [19], who gave an equivalent expression for the right-hand side of (32) and verified (34) without proving analytically that their formula corresponds to the Mellin transform of a probability distribution. The first proof of the existence of the Selberg integral distribution in full generality was given in [37], where we also discovered the decomposition in (42), followed by a new, purely probabilistic proof of (46) in [39]. The involution invariance of the Mellin transform in the equivalent form of self-duality, see (111) below, was first discovered in the special case of $\lambda_1 = \lambda_2 = 0$ in [19]. We extended it to the general case in the form of (42) in [39], followed by the general form of self-duality in [18]. The interested reader can find additional information about the Selberg integral distribution such as negative moments, asymptotic expansion, moment determinacy questions, functional equations, and infinite factorizations in [37].

3 A New Family of Barnes Beta Distributions

Let $M \in \mathbb{N}$, $a = (a_1, \cdots, a_M)$, and $b = (b_0, b_1, \cdots, b_{M-1})$, all assumed to be positive. In other words, $N = M - 1$ in the sense of Barnes beta distributions. Define

$$\eta_{M,M-1}(q|a, b) \triangleq \exp \left( (\mathcal{J}_{M-1} \log \Gamma_M)(q|a, b) - (\mathcal{J}_{M-1} \log \Gamma_M)(0|a, b) \right).$$

(43)

For example, in the case of $M = 2$ we have

$$\eta_{2,1}(q|a, b) = \frac{\Gamma_2(q + b_0|a)}{\Gamma_2(b_0|a)} \frac{\Gamma_2(b_0 + b_1|a)}{\Gamma_2(q + b_0 + b_1|a)}.$$ (44)

Such products were first discovered in [26] and [27] and then studied in depth in [28] in the context of the Mellin transform of certain functionals of the stable Lévy process.

Theorem 3.1 (Existence and Structure) Assume

$$\text{Re}(q) > -b_0.$$ (45)

$\eta_{M,M-1}(q|a, b)$ is the Mellin transform of a probability distribution $\beta_{M,M-1}(a, b)$ on $(0, \infty)$.

$$\mathbf{E} [\beta_{M,M-1}(a, b)^q] = \eta_{M,M-1}(q|a, b).$$ (46)

The distribution $\log \beta_{M,M-1}(a, b)$ is infinitely divisible and absolutely continuous on $\mathbb{R}$ and has the Lévy-Khinchine decomposition

$$\mathbf{E} \left[ \exp(q \log \beta_{M,M-1}(a, b)) \right] = \exp \left( \int_0^\infty (e^{-tq} - 1 + qt) e^{-b_0 t} e^{-\sum_{j=1}^{M-1} \sum_{i=1}^j b_i} \frac{dt}{t} \right) + q \int_0^\infty \left[ \frac{e^{-t}}{t} \frac{\prod_{i=1}^M b_i}{\prod_{i=1}^M \left( 1 - e^{-a_i t} \right)} - e^{-b_0 t} \right] \frac{dt}{t}.$$ (47)
\( \eta_{M,M-1}(q|a,b) \) satisfies the functional equation in (23) and moment formulas in (28) and (29). Given \( \kappa > 0 \), the scaling invariance property of \( \beta_{M,M-1}(a,b) \) is

\[
\beta_{M,M-1}^{\kappa} = \frac{\kappa}{\beta_{M,M-1}} \quad \text{in law,}
\]

(48)

**Theorem 3.2 (Asymptotics)** Given \( |\arg(q)| < \pi \),

\[
\eta_{M,M-1}(q|b) = \exp \left( \left( \prod_{j=1}^{M-1} b_j / \prod_{i=1}^{M} a_i \right) q \log(q) + O(q) \right), \quad q \to \infty.
\]

(49)

The Stieltjes moment problem for \( \beta_{M,M-1}(a,b) \) is determinate (unique solution) iff

\[
\prod_{j=1}^{M-1} b_j \leq 2 \prod_{i=1}^{M} a_i.
\]

(50)

It is also interesting to look at the ratio of two independent Barnes beta distributions of type \((M,M-1)\) as this ratio has remarkable factorization properties. Let

\[
\eta \triangleq (\bar{b}_0, \bar{b}_1, \cdots | a_1, \cdots | a_M).
\]

(51)

for some fixed \( \bar{b}_0 > 0 \) and define

\[
\beta_{M,M-1}(a,b,\bar{b}) \triangleq \beta_{M,M-1}(a,b) \beta_{M,M-1}^{-1}(a,\bar{b}).
\]

(52)

Then, the Mellin transform of \( \beta_{M,M-1}(a,b,\bar{b}) \) satisfies two factorizations and scaling invariance that are similar to those of Barnes beta distributions.

**Theorem 3.3 (Properties)** Let \( M \in \mathbb{N}, \bar{b}_0 > \text{Re}(q) > -b_0, i = 1 \cdots M, \Omega \triangleq \sum_{i=1}^{M} n_i a_i. \)

\[
\eta_{M,M-1}(q|a,b,\bar{b}) = \prod_{k=0}^{\infty} \eta_{M,M-1}(q+k a_i|\bar{a}_i, \bar{b}) \frac{\eta_{M,M-1}|a_i, \bar{b})}{\eta_{M,M-1}|a_i|, \bar{b})},
\]

(53)

\[
\eta_{M,M-1}(q|a,b,\bar{b}) = \prod_{n_i=0}^{\infty} \left[ \frac{b_0 + \Omega}{q + b_0 + \Omega} \frac{\bar{b}_0 + \Omega}{-q + b_0 + \Omega} \times \right.
\]

\[
\times \prod_{j=1}^{M-1} \frac{q + b_0 + b_j + \Omega}{b_0 + b_j + \Omega} \frac{-q + b_0 + b_j + \Omega}{b_0 + b_j + \Omega}
\]

\[
\times \prod_{j_i < j_2}^{M-1} \frac{b_0 + b_j + b_j + \Omega}{q + b_0 + b_j + b_j + \Omega} \frac{-q + b_0 + b_j + b_j + \Omega}{b_0 + b_j + b_j + \Omega} \cdots \right].
\]

(54)

Probalistically, these factorizations are equivalent to, respectively,

\[
\beta_{M,M-1}(a,b,\bar{b}) \text{ in law } \prod_{k=0}^{\infty} \beta_{M,M-1}(a_i, b_0 + k a_i, b_1, \cdots | b_{M-1}) \times \beta_{M,M-1}^{-1}(a_i, \bar{b}_0 + k a_i, b_1, \cdots | b_{M-1}),
\]

(55)

\[
\beta_{M,M-1}(a,b,\bar{b}) \text{ in law } \prod_{n_i=0}^{\infty} \beta_{0,M-1}(b_0 + \Omega, b_1, \cdots | b_{M-1}) \times \beta_{0,M-1}^{-1}(b_0 + \Omega, b_1, \cdots | b_{M-1}).
\]

(56)

Let \( \kappa > 0 \). Then,

\[
\beta_{M,M-1}^{\kappa}(a,b,\kappa \bar{b}) \text{ in law } \beta_{M,M-1}(a,b),
\]

(57)
Remark The multiple sine function [25] is defined by
\[ S_M(w|a) \triangleq \frac{\Gamma_M(|a| - w|a|)^{(-1)^M}}{\Gamma(w|a)}, \]  
where \( M = 0, 1, 2 \cdots \), \(|a| = \sum_{i=1}^{M} a_i \), and \( a = (a_1, \cdots, a_M) \) are fixed positive constants. Let \( \bar{b}_0 = |a| - \sum_{j=0}^{M-1} b_j > 0 \). Then, for this particular choice of \( \bar{b}_0 \) we have the interesting identity
\[ \eta_{M,M-1}(q|a,b,\bar{b}) = \exp\left( (\mathcal{S}_{M-1} \log S_M)(0|a,b) - (\mathcal{S}_{M-1} \log S_M)(q|a,b) \right), \]
\[ = \frac{S_M(b_0|a)}{S_M(q+b_0)|a|} \prod_{j=1}^{M-1} \frac{S_M(q+b_0+b_j|a)}{S_M(b_0+b_j|a)} \times \prod_{j_1<j_2} \frac{S_M(b_0+b_j_1+b_j_2|a)}{S_M(q+b_0+b_j_1+b_j_2|a)} \times \prod_{j_1<j_2<j_3} \frac{S_M(q+b_0+b_j_1+b_j_2+b_j_3|a)}{S_M(b_0+b_j_1+b_j_2+b_j_3|a)}, \]
until all \( M-1 \) indices are exhausted.

We will illustrate the general theory with the special cases of \( \beta_{1,0} \) and \( \beta_{2,1} \). Let \( M = 1 \).
\[ E[\beta_{1,0}(a,b)^q] = \frac{\Gamma_1(q+b_0|a)}{\Gamma_1(b_0|a)}, \]
\[ = a^q \frac{\Gamma\left(\frac{q+b_0}{a}\right)}{\Gamma\left(\frac{b_0}{a}\right)}, \]
by (12) so that \( \beta_{1,0} \) is a Fréchet distribution. This distribution plays an important role in the structure of both Selberg and Morris integral probability distributions. For example, the \( Y \) distribution in (36) is
\[ Y = \tau^{1/\tau} \beta_{1,0}^{-1}(a = \tau, b_0 = \tau). \]
In particular, if \( \bar{b}_0 = a - b_0 \) as in (59),
\[ E[\beta_{1,0}(a,b,\bar{b})^q] = \frac{\Gamma\left(\frac{q+b_0}{a}\right)}{\Gamma\left(\frac{b_0}{a}\right)} \frac{\Gamma\left(1 - \frac{q+b_0}{a}\right)}{\Gamma\left(1 - \frac{b_0}{a}\right)} \frac{\sin\left(\frac{\pi b_0}{a}\right)}{\sin\left(\frac{\pi a+b_0}{a}\right)}, \]
\[ = \frac{\sin\left(\frac{\pi b_0}{a}\right)}{\sin\left(\frac{\pi a+b_0}{a}\right)} \text{ for } \pi a+b_0 > 0, \]
Now, let \( M = 2, a = (a_1,a_2), b = (b_0,b_1), \bar{b} = (a_1 + a_2 - b_0 - b_1,b_1) \).
\[ E[\beta_{2,1}(a,b)^q] = \frac{\Gamma_2(q+b_0|a)}{\Gamma_2(b_0|a)} \frac{\Gamma_2(b_0+b_1|a)}{\Gamma_2(q+b_0+b_1|a)}, \]
\[ E[\beta_{2,1}(a,b,\bar{b})^q] = \frac{S_2(b_0|a)}{S_2(q+b_0|a)} \frac{S_2(q+b_0+b_1|a)}{S_2(b_0+b_1|a)} \frac{S_2(q+b_0+b_1|a)}{S_2(q+b_0+b_1|a)}, \]
For reader’s convenience, the proofs of Theorems 3.1–3.3 are deferred to Section 6.
4 Application to the Morris Integral

In this section we will consider the problem of finding a positive probability distribution having the property that its positive integral moments are given by the Morris integral. Throughout this section we let \( \tau > 1 \) and restrict \( \lambda_1, \lambda_2 \geq 0 \) for simplicity. As we did in Section 2, we write \( \tau \) as an abbreviation of \( a = (1, \tau) \) in the list of parameters of the double gamma function and \( \beta_{2,2}(a, b) \). Recall the Morris integral, see Chapter 4 of [13],

\[
\int_{-\pi}^{\pi} \prod_{l=1}^{n} e^{i\theta_l(\lambda_1 - \lambda_2)/2} |1 + e^{i\theta_l}|^{\lambda_1 + \lambda_2} \prod_{k<l} |e^{i\theta_k} - e^{i\theta_l}|^{-2/\tau} d\theta =
\]

\[
= (2\pi)^n \prod_{j=0}^{n-1} \frac{\Gamma(1 + \lambda_1 + \lambda_2 - \frac{j}{\tau}) \Gamma(1 - (j+1)/\tau)}{\Gamma(1 + \lambda_1 - \frac{j}{\tau}) \Gamma(1 + \lambda_2 - \frac{j}{\tau}) \Gamma(1 - \frac{j}{\tau})}.
\]

(68)

We wish to construct a probability distribution such that its \( n \)th moment coincides with the right-hand side of (68) for \( n < \tau \). The interest in such a construction will become clear in the next section.

Theorem 4.1 (Existence and Properties) The function

\[
\mathcal{M}(q | \tau, \lambda_1, \lambda_2) = \frac{(2\pi \tau^\frac{q}{\tau})^\tau}{\Gamma(1 - \frac{q}{\tau})} \frac{\Gamma_2(\tau(\lambda_1 + \lambda_2 + 1) + 1 - q | \tau)}{\Gamma_2(\tau(\lambda_1 + \lambda_2 + 1) + 1 | \tau)} \times
\]

\[
\times \frac{\Gamma_2(\tau(1 + \lambda_1) + 1 | \tau)}{\Gamma_2(\tau(1 + \lambda_1) + 1 - q | \tau)} \Gamma_2(\tau(1 + \lambda_2) + 1 | \tau)
\]

(69)

reproduces the product in Eq. (68) when \( q = n < \tau \) and is the Mellin transform of the distribution

\[
M_{(\tau, \lambda_1, \lambda_2)} = \frac{2\pi \tau^{1/\tau}}{\Gamma(1 - 1/\tau)} \beta_{2,2}^{-1}(\tau, b_0 = \tau, b_1 = 1 + \tau \lambda_1, b_2 = 1 + \tau \lambda_2) \times
\]

\[
\times \beta_{1,0}^{-1}(\tau, b_0 = \tau(\lambda_1 + \lambda_2 + 1) + 1),
\]

(70)

where \( \beta_{2,2}^{-1}(a, b) \) is the inverse Barnes beta of type \((2, 2)\) and \( \beta_{1,0}^{-1}(a, b) \) is the independent inverse Barnes beta of type \((1, 0)\). In particular, \( \log M_{(\tau, \lambda_1, \lambda_2)} \) is infinitely divisible, the Stieltjes moment problem for \( M_{(\tau, \lambda_1, \lambda_2)}^{-1} \) is determinate (unique solution), and the negative moments are given by

\[
E[M_{(\tau, \lambda_1, \lambda_2)}^{-n}] = (2\pi)^{-n} \prod_{j=0}^{n-1} \frac{\Gamma(1 + \lambda_1 + \frac{(j+1)}{\tau}) \Gamma(1 + \lambda_2 + \frac{(j+1)}{\tau}) \Gamma(1 - \frac{j}{\tau})}{\Gamma(1 + \lambda_1 + \lambda_2 + \frac{(j+1)}{\tau}) \Gamma(1 + \frac{j}{\tau})}.
\]

(71)

The Mellin transform is involution invariant under

\[
\tau \rightarrow \frac{1}{\tau}, \quad q \rightarrow \frac{q}{\tau}, \quad \lambda_i \rightarrow \tau \lambda_i.
\]

(72)

\[
\mathcal{M} \left( \frac{q}{\tau} \right) = \mathcal{M} \left( \frac{1}{\tau}, \tau \lambda_1, \tau \lambda_2 \right) (2\pi)^{\frac{q}{\tau}} \Gamma \left( 1 - \frac{1}{\tau} \right) \Gamma \left( 1 - \frac{q}{\tau} \right) \times
\]

\[
\times \Gamma \left( 1 - \frac{1}{\tau} \right) \Gamma \left( 1 - \frac{q}{\tau} \right).
\]

(73)
In the special case of $\lambda_1 = \lambda_2 = 0$, we have

$$
\Omega(q | \tau, 0, 0) = \frac{(2\pi)^q}{\Gamma(1 - \frac{q}{\tau})} \Gamma(1 - \frac{q}{\tau}),
$$

(74)

$$
M(\tau, 0, 0) = \frac{2\pi \tau^{1/\tau}}{\Gamma(1 - 1/\tau)} \beta^{-1}_{1,0}(\tau, b_0 = \tau).
$$

(75)

**Remark** The special case of $\lambda_1 = \lambda_2 = 0$ was first treated in [14] and corresponds to the Dyson integral. The involution invariance in this case, in its equivalent self-duality form, see (111) below, was first discovered in [19].

The proofs are deferred to Section 6.

## 5 Maximum Distribution of the Gaussian Free Field on the Unit Interval and Circle

In this section we will formulate precise conjectures about the distribution of the maximum of the discrete 2D gaussian free field (GFF) on the unit interval and circle with a non-random logarithmic potential. We will not attempt to review the GFF here but rather refer the reader to [14] for the circle case and to [19] and Section 3 of [40] for the interval case. Suffice it to say that our approach to the GFF construction is essentially based on the construction of [1], [2], [34].

**GFF on the interval** Let $V_\epsilon(x)$ be a centered gaussian process on $[0, 1]$ with the covariance

$$
\text{Cov}[V_\epsilon(u), V_\epsilon(v)] = \begin{cases} 
-2 \log |u - v|, & \epsilon \leq |u - v| \leq 1, \\
2 \left(1 - \log \epsilon - \frac{|u - v|}{\epsilon}\right), & |u - v| \leq \epsilon.
\end{cases}
$$

(76)

The limit $\lim_{\epsilon \to 0} V_\epsilon(x)$ is called the continuous GFF on the interval $x \in [0, 1]$ and its discretized version $V_\epsilon(x_i), x_i = i\epsilon, i = 0 \cdots N, \epsilon = 1/N$ is the discrete GFF on the interval.

We note that in applications, see [40] for example, the GFF construction arises in a slightly more general form of

$$
\text{Cov}[V_\epsilon(u), V_\epsilon(v)] = \begin{cases} 
-2 \log |u - v|, & \epsilon \ll |u - v| \leq 1, \\
2 (\kappa - \log \epsilon), & u = v, \\
+ O(\epsilon), & \text{else},
\end{cases}
$$

(77)

where $\kappa \geq 0$ is some fixed constant and the details of regularization are relegated to the $O(\epsilon)$ term. It is worth emphasizing that the choice of covariance regularization for $|u - v| \leq \epsilon$ has no effect on distribution of the maximum, so long as the variance behaves as in (77), due to Theorem 6 in [2] as explained below. The same remark applies to the GFF on the circle so we give its definition in the more general form.
**GFF on the circle** Let \( V_\varepsilon(\psi) \) be a centered gaussian process with the covariance

\[
\text{Cov} \left[ V_\varepsilon(\psi), V_\varepsilon(\xi) \right] = \begin{cases} 
-2 \log |e^{i\psi} - e^{i\xi}|, & |\xi - \psi| \gg \varepsilon, \\
2 (\kappa - \log \varepsilon), & \psi = \xi,
\end{cases}
\]

\[+ O(\varepsilon), \quad (78)\]

where \( \kappa \geq 0 \) is some fixed constant. The limit \( \lim_{\varepsilon \to 0} V_\varepsilon(\psi) \) is called the GFF on the circle \( \psi \in [-\pi, \pi) \) and its discretized version \( V_\varepsilon(\psi_j), \psi_j = j\varepsilon, j = -N/2 \cdots N/2, \varepsilon = 1/N \) is the discrete GFF on the circle.

The existence of such objects follows from the general theory of [41] as shown in [2] and [34].

**Problem formulation for the interval** Let \( \lambda_1, \lambda_2 \geq 0 \). Let

\[ N = 1/\varepsilon \quad (79)\]

so we can think of \( V_\varepsilon(x) \) as being defined at \( x_i = i\varepsilon, i = 0 \cdots N \). What is the distribution of

\[ V_N = \max \{ V_\varepsilon(x_i) + \lambda_1 \log(x_i) + \lambda_2 \log(1 - x_i), i = 1 \cdots N \} \quad (80)\]

in the form of an asymptotic expansion in \( N \) in the limit \( N \to \infty \)?

**Problem formulation for the circle** Let \( \alpha \geq 0 \). Let

\[ N = 2\pi/\varepsilon \quad (81)\]

and think of \( V_\varepsilon(\psi) \) as being defined at \( \psi_j = j\varepsilon, j = -N/2 \cdots N/2 \). What is the distribution of

\[ V_N = \max \{ V_\varepsilon(\psi_j) + 2\alpha \log |1 + e^{i\psi_j}|, j = -N/2 \cdots N/2 \} \quad (82)\]

in the form of an asymptotic expansion in \( N \) in the limit \( N \to \infty \)?

We will consider the case of the GFF on the interval first. Define the critical Selberg integral probability distribution in terms of the Selberg integral distribution in Theorem 2.6 by the formula

**Critical Selberg Integral Distribution**

\[ M_{(\tau=1,\lambda_1,\lambda_2)} \triangleq \lim_{\tau \downarrow 1} \frac{\Gamma \left( 1 - 1/\tau \right)}{2\pi} M_{(\tau,\lambda_1,\lambda_2)} \quad (83)\]

This limit exists by (40). We will refer to \( M_{(\tau=1,\lambda_1,\lambda_2)} \) as the critical Selberg integral distribution.

**Conjecture 5.1 (Maximum of the GFF on the Interval)** The leading asymptotic term in the expansion of the Laplace transform of \( V_N \) in \( N \) is

\[ E[e^{qV_N}] \approx e^{q(2\log N - (3/2)\log \log N + \text{const})} E[ M_{(\tau=1,\lambda_1,\lambda_2)}^q ], \quad N \to \infty. \quad (84)\]
Probabilistically, let

\[ X_1 \triangleq \beta_{2,2}^{-1}(\tau = 1, b_0 = 2 + \lambda_1, b_1 = (\lambda_2 - \lambda_1)/2, b_2 = (\lambda_2 - \lambda_1)/2), \]
\[ X_2 \triangleq \beta_{2,2}^{-1}(\tau = 1, b_0 = 2 + (\lambda_1 + \lambda_2)/2, b_1 = 1/2, b_2 = 1/2), \]
\[ X_3 \triangleq \beta_{2,2}^{-1}(\tau = 1, b_0 = 2, b_1 = (2 + \lambda_1 + \lambda_2)/2, b_2 = (2 + \lambda_1 + \lambda_2)/2), \]

be as in Theorem 2.6 with \( \tau = 1 \). Then, as \( N \to \infty \),

\[ V_N = 2\log N - \frac{3}{2} \log \log N + \text{const} + \mathcal{N}(0, 4\log 2) + \log X_1 + \log X_2 + \log X_3 + \]

\[ + \log Y + \log Y' + o(1), \quad (89) \]

where \( Y' \) is an independent copy of \( Y \).

**Remark** This conjecture at the level of the Mellin transform is due to [19] in the case of \( \lambda_1 = \lambda_2 = 0 \) and [18] for general \( \lambda_1 \) and \( \lambda_2 \). Our probabilistic re-formulation of their conjecture is new. We also note that the Mellin transform of the Barnes beta distribution \( \beta_{2,2} \) with \( a = (1, 1) \) can be expressed in terms of the Barnes \( G \) function, see [3], as follows.

\[
E[\beta_{2,2}^q(1, b_0, b_1, b_2)] = \frac{G(b_0) G(q + b_0 + b_1) G(q + b_0 + b_2)}{G(q + b_0) G(b_0 + b_1) G(b_0 + b_2)} \times
\]

\[ \times \frac{G(b_0 + b_1 + b_2)}{G(q + b_0 + b_1 + b_2)}. \quad (90) \]

In the special case of \( \lambda_1 = \lambda_2 = 0 \), \( X_1 \) degenerates, \( X_3 \) becomes Pareto, and

\[ X_2 = \beta_{2,2}^{-1}(1, b_0 = 2, b_1 = 1/2, b_2 = 1/2), \quad (91) \]

see Section 6 of [39] for details.

Similarly, to formulate our conjecture for the maximum of the GFF on the circle, we need to define the critical Morris integral probability distribution.

**Critical Morris Integral Distribution**

\[ M_{(\tau=1,\lambda_1,\lambda_2)} \triangleq \lim_{\tau \uparrow 1} \frac{\Gamma(1 - 1/\tau)}{2\pi} M_{(\tau,\lambda_1,\lambda_2)}, \quad (92) \]

where \( M_{(\tau,\lambda_1,\lambda_2)} \) is as in Theorem 4.1.

This limit exists by (70). We will refer to \( M_{(\tau=1,\lambda_1,\lambda_2)} \) as the critical Morris integral distribution.

**Conjecture 5.2 (Maximum of the GFF on the Circle)** The leading asymptotic term in the expansion of the Laplace transform of \( V_N \) in \( N \) is

\[ E[e^{tV_N}] \approx e^{t(2\log N - (3/2)\log \log N + \text{const})} E[M_{(\tau=1,\alpha,\alpha)}^q], \quad N \to \infty, \quad (93) \]
Probabilistically, let
\[ X ≜ \beta_{2,2}^{-1}(\tau = 1, b_0 = 1, b_1 = 1 + \alpha, b_2 = 1 + \alpha), \]  
\[ Y ≜ \beta_{1,0}^{-1}(\tau = 1, b_0 = 2\alpha + 2), \]  
\[ Y' ≜ \beta_{1,0}^{-1}(\tau = 1, b_0 = 1) \]  
be as in Theorem 4.1 with \( \tau = 1 \). Then,
\[ V_N = 2 \log N - \frac{3}{2} \log \log N + \text{const} + \log X + \log Y + \log Y' + o(1). \]  
\[ (97) \]
If \( \alpha = 0 \),
\[ V_N = 2 \log N - \frac{3}{2} \log \log N + \text{const} + \log Y + \log Y' + o(1), \]  
\[ (98) \]
where
\[ Y \overset{\text{inlaw}}{=} Y' = \beta_{1,0}^{-1}(\tau = 1, b_0 = 1). \]  
\[ (99) \]
Remark This conjecture is due to [14] in the case of \( \alpha = 0 \). The extension to general \( \alpha \) is new.

In the rest of this section we will give a heuristic derivation of our conjectures. Our approach is based on the freezing hypothesis and calculations of [14] and [19] as well as our conjectures relating the distribution of the exponential functional of the GFF on the interval and circle to the Selberg and Morris integral distributions, respectively. To this end, we need to briefly recall the principal result of [2] concerning the existence of the gaussian multiplicative chaos measure (also known as the limit lognormal measure, see [32], [33]), which we are stating here in slightly more general form.

**Theorem 5.3 (Multiplicative Chaos on the Interval)** Let \( 0 \leq \beta < 1 \). The exponential functional of the GFF on the interval converges weakly a.s. as \( \varepsilon \to 0 \) to a non-degenerate limit random measure
\[ e^{-\beta^2(x - \log \varepsilon)} \int_a^b e^{BV_\varepsilon(u)} du \to M_\beta(a, b), \]  
\[ (100) \]
\[ \mathbb{E}[M_\beta(a, b)] = |b - a|. \]  
\[ (101) \]
The moments of the (generalized) total mass of the limit measure are given by the Selberg integral: let \( n < 1/\beta^2 \),
\[ \mathbb{E} \left[ \left( \int_0^1 s^{\lambda_1}(1 - s)^{\lambda_2} M_\beta(ds) \right)^n \right] = \int \prod_{i=1}^n s_i^{\lambda_1}(1 - s_i)^{\lambda_2} \prod_{i<j} |s_i - s_j|^{-2\beta^2} ds_1 \cdots ds_n. \]  
\[ (102) \]
We constructed a probability distribution having the required moments, see Theorem 2.6, which leads us to the following conjecture, originally due to [36] for \( \lambda_1 = \lambda_2 = 0 \) and due to [19] (at the level of the Mellin transform, see Remark 2 at the end of Section 2) and [37] in general.

**Conjecture 5.4 (Law of Total Mass)** Let \( M_{(\tau, \lambda_1, \lambda_2)} \) be the Selberg integral distribution in (33).
\[ M_{(\tau, \lambda_1, \lambda_2)} \overset{\text{inlaw}}{=} \int_0^1 s^{\lambda_1}(1 - s)^{\lambda_2} M_\beta(ds), \tau = 1/\beta^2 > 1. \]  
\[ (103) \]
Now, following [19], define the exponential functional

$$Z_{\lambda_1, \lambda_2, \epsilon}(\beta) \triangleq \sum_{i=1}^{N} \beta \lambda_1 (1 - x_i) \beta \lambda_2 e^{\beta V_{\epsilon}(x_i)}. \quad (104)$$

Using the identity

$$P(V_N < s) = \lim_{\beta \to \infty} E\left[\exp\left(-e^{-\beta s} Z_{\lambda_1, \lambda_2, \epsilon}(\beta)/C\right)\right], \quad (105)$$

which is applicable to any sequence of random variables and an arbitrary $\beta$-independent constant $C$, the distribution of the maximum is reduced to the Laplace transform of the exponential functional in the limit $\beta \to \infty$. Now, by Theorem 5.3, it is known that for $0 < \beta < 1$ the exponential functional converges as $N \to \infty$ to the (generalized) total mass of the multiplicative chaos (limit lognormal) measure on the unit interval, which is conjectured to be given by the Selberg integral distribution, resulting in the approximation

$$Z_{\lambda_1, \lambda_2, \epsilon}(\beta) \approx N^{1+\beta^2} e^{\beta^2 \kappa} M_{(\tau, \lambda_1, \lambda_2)}, \quad N \to \infty, \quad (106)$$

where

$$\tau = 1/\beta^2, \quad 0 < \beta < 1. \quad (107)$$

Next, we recall the involution invariance of the Mellin transform of the Selberg integral distribution, see (42). Denoting the Mellin transform as in Theorem 2.6,

$$M(q \mid \tau, \lambda_1, \lambda_2) = E\left[M_{(\tau, \lambda_1, \lambda_2)}^q\right], \quad \text{Re}(q) < \tau, \quad (108)$$

we have

$$M\left(\frac{q}{\tau}, 1/\tau, \tau, \lambda_1, \lambda_2 \right)(2\pi)^{-\frac{q}{\tau}} \Gamma^\frac{q}{\tau}(1 - \frac{q}{\tau}) \Gamma(1 - \frac{q}{\tau}) \approx M(q \mid \tau, \lambda_1, \lambda_2)(2\pi)^{-q} \times \Gamma^q(1 - \frac{1}{\tau}) \Gamma(1 - q). \quad (109)$$

Hence, introducing the function

$$F(q \mid \beta, \lambda_1, \lambda_2) \triangleq M\left(\frac{q}{\beta}, 1/\beta, \beta, \lambda_1, \lambda_2 \right)(2\pi)^{-\frac{q}{\beta}} \Gamma^\frac{q}{\beta}(1 - \frac{q}{\beta}) \Gamma(1 - \frac{q}{\beta}), \quad (110)$$

one observes that it satisfies the identity

$$F(q \mid \beta, \lambda_1, \lambda_2) = F\left(q \mid \frac{1}{\beta}, \lambda_1, \lambda_2 \right), \quad (111)$$

first discovered in [19] in the special case of $\lambda_1 = \lambda_2 = 0$, then formulated in general in the form of (109) in [39], and in [18] in this form. On the other hand, as shown in [19], one has the general identity

$$\int_{\mathbb{R}} e^{q y} \frac{d}{dy} \exp\left(-e^{-\beta y} X\right) dy = X^\frac{q}{\beta} \Gamma\left(1 - \frac{q}{\beta}\right), \quad \text{Re}(q) < 0, \quad X > 0. \quad (112)$$

2 Theorem 6 in [2] shows that the laws of the total mass of the continuous and discrete multiplicative chaos measures are the same provided the $\epsilon$ parameter in (76) coincides with the discretization step.

3 The validity of approximating the finite $N$ quantity with the $N \to \infty$ limit is discussed in [14].
Letting
\[ X = Z_{\lambda_1, \lambda_2, e}(\beta) \frac{e^{\kappa} \Gamma(1 - \beta^2)}{2\pi} \] (113)
onone sees by means of (106) that for \( 0 < \beta < 1 \),
\[ E[X^\beta] \approx e^{q\kappa(\beta + \frac{1}{\beta})} \mathcal{N}(\beta + \frac{1}{\beta}) \mathcal{M}(\frac{q}{\beta}, \beta \lambda_1, \beta \lambda_2) \frac{1}{(2\pi)^{\frac{3}{2}} \Gamma(1 - \beta^2)} \], \( N \to \infty \), (114)
so that
\[ E[X^\beta] \Gamma(1 - \frac{q}{\beta}) \approx e^{q\kappa(\beta + \frac{1}{\beta})} \mathcal{N}(\beta + \frac{1}{\beta}) F(q | \beta, \lambda_1, \lambda_2), \] \( N \to \infty \). (115)

On the other hand, by letting the constant \( C \) in (105) to be taken to be
\[ C = \frac{2\pi}{e^{\kappa} \Gamma(1 - \beta^2)}, \] (116)
and combining (105) with (112), one obtains
\[ E[e^{qV_X}] = \lim_{\beta \to \infty} \left[ E[X^\beta] \Gamma(1 - \frac{q}{\beta}) \right]. \] (117)
The right-hand side of this equation has only been determined for \( 0 < \beta < 1 \), see (115). Due to the self-duality of the right-hand side, one assumes that it gets frozen at \( \beta = 1 \), as first formulated in [19].

**Conjecture 5.5 (The Freezing Hypothesis)** Let \( \beta > 1 \).
\[ E[X^\beta] \Gamma(1 - \frac{q}{\beta}) = E[X^\beta] \Gamma(1 - \frac{q}{\beta}) \big|_{\beta = 1}. \] (118)

One must note however that \( C \) is not \( \beta \)-independent. It is argued in [15] and [20] that the \( \Gamma(1 - \beta^2) \) term shifts the maximum by \( -(3/2) \log \log N \), see also [10]. Overall, we then obtain by (115),
\[ E[e^{qV_X}] \approx e^{(2 \log N - (3/2) \log \log N + \text{const})} F(q | \beta = 1, \lambda_1, \lambda_2), \] \( N \to \infty \) (119)
for some constant\(^4\). Finally, recalling definitions of the critical Selberg integral distribution in (83) and of \( F(q | \beta, \lambda_1, \lambda_2) \) in (110), and appropriately adjusting the constant,
\[ E[e^{qV_X}] \approx e^{(2 \log N - (3/2) \log \log N + \text{const})} M_{\tau = 1, \lambda_1, \lambda_2} \Gamma(1 - q), \] (120)
so that \( Y' \) comes from the \( \Gamma(1 - q) \) factor and has the same law as \( Y \).

The argument for the GFF on the circle goes through verbatim so we will only point out the key steps and omit redundant details. Define the exponential functional
\[ Z_{\alpha, e}(\beta) = \sum_{j = -N/2}^{N/2} |1 + e^{i \psi_j}|^{2\alpha \beta} e^{B_{\psi_j}(\psi_j)}. \] (121)
To describe its limit as \( N \to \infty \) we need to compute the distribution of the generalized total mass of the limit lognormal measure on the circle. The following result follows from the general theory of multiplicative chaos measures of [24].

\(^4\)As remarked in [21], this procedure only determines the distribution of the maximum up to a constant term.
Theorem 5.6 (Multiplicative Chaos on the Circle) Let $0 \leq \beta < 1$. The exponential functional of the GFF on the circle converges weakly a.s. as $\varepsilon \to 0$ to a non-degenerate limit random measure

$$e^{-\beta^2 (\kappa - \log \varepsilon)} \int_\phi^\Psi e^{\beta V_c(\theta)} d\theta \longrightarrow M_\beta(\phi, \Psi),$$

(122)

$$E[M_\beta(\phi, \Psi)] = |\Psi - \phi|.$$  

(123)

The moments of the (generalized) total mass of the limit measure are given by the Morris integral: let $n < 1/\beta^2$,

$$E\left(\left(\int_{-\pi}^{\pi} e^{i\int_{-\pi}^{\pi} \frac{\lambda_1 - \lambda_2}{2} dM_\beta(\psi)} + e^{i\int_{-\pi}^{\pi} \psi_1 + \lambda_2} dM_\beta(\psi)\right)^n\right) = \int_{-\pi}^{\pi} \prod_{l=1}^{n} e^{i\theta_l \lambda_1 - \lambda_2} |1 + e^{i\theta_l \lambda_1 + \lambda_2} | \times \prod_{k<l} |e^{i\theta_k} - e^{i\theta_l}|^{-2\beta^2} d\theta.$$  

(124)

We constructed a probability distribution having the required moments, see Theorem 4.1, which leads us to the following conjecture.

Conjecture 5.7 (Law of Total Mass) Let $M_{(\tau, \lambda_1, \lambda_2)}$ be as in (70). Let

$$\lambda_1 = \lambda_2 = \alpha \frac{\kappa}{2}$$

(125)

$$M_{(\tau, \alpha, \alpha)} \overset{\text{in law}}{=} \int_{-\pi}^{\pi} e^{\int_{-\pi}^{\pi} \frac{\lambda_1 - \lambda_2}{2} dM_\beta(\psi)} \frac{1 + e^{i\int_{-\pi}^{\pi} \psi_1 + \lambda_2} dM_\beta(\psi)}{\tau} = 1/\beta^2 > 1.$$  

(126)

This conjecture for $\lambda_1 = \lambda_2 = 0$ is due to [14], the general case is original to this paper.

It follows that the exponential functional in (121) can be approximated by

$$Z_{\alpha, \kappa}(\beta) \approx \left(\frac{N}{2\pi}\right)^{1+\beta^2} e^{\beta^2 \kappa} M_{(\tau, \beta, \alpha, \alpha)}, \; N \to \infty.$$  

(127)

The rest of the argument is the same as for the GFF on the interval, with (73) replacing (42).

6 Proofs

We begin with a key lemma on the action of $\mathcal{M}_{M-1}$ on polynomials.

Lemma 6.1 (Main lemma) Let $f(t)$ be a fixed function and define the associated Bernoulli polynomials by

$$B_k^{(f)}(x) \triangleq \frac{d^m}{dt^m} f(t) e^{-xt}.$$  

(128)

$$\left(\mathcal{M}_{M-1} B_k^{(f)}\right)(q \mid b) = 0, \; k < M - 1,$$

(129)

$$\left(\mathcal{M}_{M-1} B^{(f)}_{M-1}\right)(q \mid b) = \left(\mathcal{M}_{M-1} B^{(f)}_{M-1}\right)(0 \mid b),$$

(130)

$$\left(\mathcal{M}_{M-1} B^{(f)}_M\right)(q \mid b) - \left(\mathcal{M}_{M-1} B^{(f)}_M\right)(0 \mid b) = -q f(0) M! \prod_{j=1}^{M-1} b_j.$$  

(131)

This restriction is necessary as $M_{(\tau, \lambda_1, \lambda_2)}$ is real-valued whereas $\int_{-\pi}^{\pi} e^{i\int_{-\pi}^{\pi} \frac{\lambda_1 - \lambda_2}{2} dM_\beta(\psi)}$ is not in general, unless $\lambda_1 = \lambda_2$. The problem of determining the law of $\int_{-\pi}^{\pi} e^{i\int_{-\pi}^{\pi} \frac{\lambda_1 - \lambda_2}{2} dM_\beta(\psi)}$ for $\lambda_1 \neq \lambda_2$ is left to future research.
Proof The identities in (129) and (130) were first established in [38]. As the argument for (131) is similar, we will reproduce the original argument here for completeness. Define the function $g(t)$,

$$
g(t) \equiv f(t)e^{-qt} e^{-b_0 t} \prod_{j=1}^{N} (1 - e^{-b_j t}). \tag{132}
$$

Then, using the identity

$$
e^{-(q+b_0) t} \prod_{j=1}^{N} (1 - e^{-b_j t}) = e^{-(q+b_0) t} \sum_{p=0}^{N} (-1)^p \sum_{k_1 < \cdots < k_p = 1} \exp\left(- (b_{k_1} + \cdots + b_{k_p}) t\right),
$$

one obtains

$$
g^{(k)}(0) = (\mathcal{S}_N e^{-Mt})(q|b),
$$

$$
= 0, \text{ if } k < N,
$$

$$
= f(0) N! \prod_{j=1}^{N} b_j, \text{ if } k = N. \tag{135}
$$

To verify (131), letting $N = M - 1$, $k = M$, one observes that (134) implies the identity

$$(\mathcal{S}_{M-1} B_{M-1}^{(f)}(x))(q|b) = f(0) \frac{d^M}{dt^M} \left|_{t=0} \right. \prod_{j=1}^{M-1} (1 - e^{-b_j t}) + M! \prod_{j=1}^{M-1} b_j B_1^{(f)}(q+b_0). \tag{137}
$$

Hence,

$$(\mathcal{S}_{M-1} B_{M-1}^{(f)}(x))(q|b) - (\mathcal{S}_{M-1} B_{M-1}^{(f)}(x))(0|b) =
M! \prod_{j=1}^{M-1} b_j \left( B_1^{(f)}(q+b_0) - B_1^{(f)}(b_0) \right). \tag{138}
$$

Finally, it remains to notice that $B_1^{(f)}(x)$ satisfies the identity

$$B_1^{(f)}(x) - B_1^{(f)}(x+y) = f(0) y, \tag{139}
$$

and (131) follows.

Proof of Theorem 3.1 We wish to establish the identity

$$e^{-tq} - 1) e^{-b_0 t} \prod_{j=1}^{M} (1 - e^{-b_j t}) + q e^{-t} \prod_{i=1}^{M} a_i \right). \tag{140}
$$

Once it is established, (47) follows immediately by adding and subtracting $qt$ in the integrand, which allows us to split the integral in (140) into two individual integrals and thereby bring it to the required
Lévy-Khinchine form. To verify (140) we need to recall the Ruijsenaars formula in Theorem 2.1. We see from Lemma 6.1 that for any \( k = 0 \cdots M - 1 \) we have the identity

\[
(\mathcal{L}_{M-1}B_{M,k}(x|a))(q|b) - (\mathcal{L}_{M-1}B_{M,k}(x|a))(0|b) = 0.
\]

(141)

It follows upon substituting (7) into (43) and using (131) that the only non-vanishing terms are

\[
\eta_{M,M-1}(q|a,b) = \exp\left(\int_0^\infty \frac{dt}{t} \left[ \frac{(\mathcal{L}_{M-1}e^{-at})(q|b) - (\mathcal{L}_{M-1}e^{-at})(0|b)}{\prod_{i=1}^M (1 - e^{-a_i t})} \right] + \frac{M-1}{M} b_j \prod_{i=1}^M a_i \right),
\]

(142)

and the result follows from (133). Now, we have the obvious identities

\[
\int_0^\infty \frac{e^{-bt}}{t} \frac{1}{\prod_{i=1}^M (1 - e^{-a_i t})} dt = \infty,
\]

(143)

\[
\int_0^\infty e^{-bt} \frac{1}{\prod_{i=1}^M (1 - e^{-a_i t})} dt = \infty,
\]

(144)

which imply that \( \log \beta_{M,M-1}(a,b) \) is absolutely continuous and supported on \( \mathbb{R} \) by Theorem 4.23 and Proposition 8.2 in Chapter 4 of [47], respectively. The scaling invariance in (48) is a corollary of (14) and (131).

**Proof of Theorem 3.2** The asymptotic expansion of \( \log \Gamma_M(w|a) \) in (8) consists of the leading term \( -B_{M,M}(w|a) \log(\omega)/M! \) plus a polynomial remainder. Lemma 6.1 shows that the remainder term contributes \( O(q) \) to \( \log \eta_{M,M-1}(q|a,b) \). It remains to show that as \( q \to \infty \),

\[
(\mathcal{L}_{M-1}B_{M,M}(x|a) \log(x))(q|b) = -M! \left( \prod_{j=1}^{M-1} b_j / \prod_{i=1}^M a_i \right) q \log(q) + O(q).
\]

(145)

Slightly generalizing the calculation in Lemma 6.1 let

\[
g(t) = f(t)e^{-qt} \frac{d^n}{dt^n} \left[ e^{-bt} \prod_{j=1}^{M-1} (1 - e^{-b_j t}) \right].
\]

(146)

Then,

\[
g^{(n)}(0) = \sum_{m=0}^n \binom{n}{m} B_{n-m}^{(f)}(q) \frac{d^{m+r}}{dt^{m+r}}|_{t=0} \left[ e^{-bt} \prod_{j=1}^{M-1} (1 - e^{-b_j t}) \right],
\]

(147)

\[
= (-1)^r \sum_{p=0}^{M-1} (-1)^p \sum_{k_1 < \cdots < k_p} (b_0 + \sum b_{k_j})^r B_n^{(f)}(q + b_0 + \sum b_{k_j}).
\]

(148)
In our case $n = M$, $f(t)$ as in (5), and the expression in (148) is the coefficient of $1/q^r$ that one gets by expanding $(\mathcal{M}^{-1}B_{M,M}(x|a) \log(x))(q|b)$ in powers of $1/q$. By (147) we can restrict ourselves to $m + r \geq M - 1$. On the other hand, the overall power of $q$ is $M - m - r$ so that we are only interested in $M - m - r \geq 1$, as the other terms contribute $O(1)$. Hence, $m = M - r - 1$, and the contribution of such terms is $O(q)$,\[
(\mathcal{M}^{-1}B_{M,M}(x|a) \log(x))(q|b) = \log(q) (\mathcal{M}^{-1}B_{M,M}(x|a))(q|b) + O(q), \tag{149}\]
and the result follows from (131). Finally, the asymptotic behavior in (49) coincides with the asymptotic behavior of generalized gamma distributions and (50) follows from the known solution to the Stieltjes moment problem for these distributions, see [29]. \hfill \Box

**Proof of Theorem 3.3** The factorizations in (53) and (54) are corollaries of Lemma 6.1 and Shintani and Barnes factorizations of the multiple gamma functions, see (16) and (3), respectively. A direct proof can be given as follows. The Mellin transform of $\beta_{M,M-1}(a, b, \bar{b})$ is\[
\eta_{M,M-1}(q|a, b, \bar{b}) = \exp \left( \int_0^\infty \frac{dt}{t} \left( (e^{-tq}-1)e^{-b_0t} + (e^{tq}-1)e^{-b_bt} \right) \prod_{j=M}^{M-1} \left( 1 - e^{-\bar{b}_jt} \right) \right). \tag{150}\]
Let $i = M$ without any loss of generality. The formula in (150) can be written in the form\[
\eta_{M,M-1}(q|a, b, \bar{b}) = e^{j} \sum_{k=0}^{\infty} \left[ (e^{-tq}-1)e^{-(b_0+ka)t} + (e^{tq}-1)e^{-(\bar{b}_0+ka)t} \right] \prod_{j=1}^{M-1} \left( 1 - e^{-b_jt} \right) \times \exp \left( \int_0^\infty \frac{dt}{t} (e^{tq}-1)e^{-(\bar{b}_0+ka)t} \prod_{j=1}^{M-1} \left( 1 - e^{-\bar{b}_jt} \right) \right), \tag{151}\]
which is exactly the same as the expression in (53) if one recalls (21) and the following identity that we first noted in (38),\[
\eta_{M,N}(q|a, b_0 + x, b_1 \cdots b_N) \eta_{M,N}(x|a, b) = \eta_{M,N}(q+x|a, b), \quad x > 0, \tag{152}\]
where $b = (b_0, b_1, \cdots b_N)$. The expression in (151) is equivalent to (55). If we now apply (27) to each Barnes beta factor in (55), we obtain (56), which is equivalent to (54). The scaling invariance in (57) follows from (48). \hfill \Box

**Proof of Theorem 4.1** Recall the functional equation of the double gamma function, see (11) with $M = 2$, and the definition of $\Gamma_1(w|a)$ in (12). Using the functional equation repeatedly, we obtain for $a = (1, \tau)$ and $k \in \mathbb{N}$,\[
\frac{\Gamma_2(z+1-k|\tau)}{\Gamma_2(z+1|\tau)} = \left( \frac{1}{2\pi \tau} \right)^{k/2} \sum_{j=0}^{k-1} (z-j)/\tau \prod_{j=0}^{k-1} \Gamma \left( \frac{z-j}{\tau} \right). \tag{153}\]
We now apply this equation to each of the four ratios of double gamma functions in (69) with \( q = n \), which results in (68). Now, the inverse Barnes beta distribution \( \beta^{-1}_{2,2}(a, b) \) with parameters \( a = (1, \tau) \) and \( b = (\tau, 1 + \tau \lambda_1, 1 + \tau \lambda_2) \) has the Mellin transform

\[
\mathbb{E}\left[ \beta^{-q}_{2,2}(\tau, 1 + \tau \lambda_1, 1 + \tau \lambda_2) \right] = \frac{\Gamma_2(-q + \tau | \tau)}{\Gamma_2(\tau | \tau)} \frac{\Gamma_2(\tau(1 + \lambda_1) + 1 | \tau)}{\Gamma_2(\tau(1 + \lambda_1) + 1 - q | \tau)} \times \frac{\Gamma_2(\tau(1 + \lambda_2) + 1 | \tau)}{\Gamma_2(\tau(1 + \lambda_2) + 1 - q | \tau)} \times \frac{\Gamma_2(\tau(\lambda_1 + \lambda_2 + 1) + 2 - q | \tau)}{\Gamma_2(\tau(\lambda_1 + \lambda_2 + 2) + 2 | \tau)}. \tag{154}
\]

The difference between this expression and that in (69) is in the last factor. Applying the functional equation once again, we get

\[
\frac{\Gamma_2(\tau(\lambda_4 + \lambda_2 + 1) + 1 - q | \tau)}{\Gamma_2(\tau(\lambda_1 + \lambda_2 + 1) + 1 | \tau)} = \tau^{-\frac{\lambda_1 + \lambda_2 + 1 + \frac{1 - q}{\tau}}{\Gamma(\lambda_1 + \lambda_2 + 1 + \frac{1}{\tau})}} \times \frac{\Gamma_2(\tau(\lambda_4 + \lambda_2 + 1) + 2 - q | \tau)}{\Gamma_2(\tau(\lambda_1 + \lambda_2 + 2) + 2 | \tau)}. \tag{155}
\]

Recalling the definition of the Mellin transform of the inverse Barnes beta \( \beta^{-1}_{1,0}(a, b) \) with parameters \( a = \tau \) and \( b = 1 + \tau(1 + \lambda_1 + \lambda_2) \),

\[
\mathbb{E}\left[ \beta^{-q}_{1,0}(\tau, 1 + \tau(1 + \lambda_1 + \lambda_2)) \right] = \tau^{-\frac{\lambda_1 + \lambda_2 + 1 + \frac{1 - q}{\tau}}{\Gamma(1 + \lambda_1 + \lambda_2 + 1 + \frac{1}{\tau})}}. \tag{156}
\]

we see that the Mellin transform of the distribution \( M_{(\tau, \lambda_1, \lambda_2)} \) in (70) coincides with the expression in (69).

The infinite divisibility of \( \log M_{(\tau, \lambda_1, \lambda_2)} \) follows from that of \( \log \beta^{-1}_{2,2}(a, b) \) and \( \log \beta^{-1}_{1,0}(a, b) \). The determinacy of the Stieltjes moment problem for \( M^{-1}_{(\tau, \lambda_1, \lambda_2)} \) follows from Theorem 3.2 as \( \beta_{2,2}(a, b) \) is compactly supported. In our case the condition for determinacy in (50),

\[
1 \leq 2 \tau, \tag{157}
\]

is satisfied as \( \tau > 1 \). The negative moments and the special case of \( \lambda_1 = \lambda_2 = 0 \) follow from (153). To prove the involution invariance in (73), we need to recall the scaling property of the multiple gamma function, see (14). The transformation in (72) corresponds to \( \kappa = 1/\tau \). We note first that

\[
\frac{1}{\tau}(1, \tau) = (1, \frac{1}{\tau}) \tag{158}
\]

in the sense of the parameters of the double gamma function. We apply (14) to each of the double gamma factors in (69) under the transformation in (72). For example,

\[
\Gamma_2\left(\frac{1}{\tau}(\tau \lambda_1 + \tau \lambda_2 + 1) + 1 - q \left| \frac{1}{\tau}\right.\right) = \Gamma_2\left(\frac{1}{\tau}(\tau(\lambda_1 + \lambda_2 + 1) + 1 - q) \left| \frac{1}{\tau}\right.\right) = \left(\frac{1}{\tau}\right)^{-B_{2,2}(\tau(\lambda_1 + \lambda_2 + 1) + 1 - q)/2} \times \Gamma_2\left(\tau(\lambda_1 + \lambda_2 + 1) + 1 - q \left| \tau\right.\right). \tag{159}
\]
Using the identity
\[ B_{2,2}(x|a) = \frac{x^2}{a_1a_2} - \frac{x(a_1 + a_2)}{a_1a_2} + \frac{a_1^2 + 3a_1a_2 + a_2^2}{6a_1a_2}, \]  
with \((a_1 = 1, a_2 = \tau)\), we collect all terms and simplify to obtain
\[ M\left(\frac{q}{\tau}, \tau \lambda_1, \tau \lambda_2\right)(2\pi)^{-\frac{q}{2}} \Gamma^q(1 - \tau) = M(q|\tau, \lambda_1, \lambda_2) \left(\frac{2\pi \tau^q}{\Gamma(1 - \frac{q}{2})}\right)^{-q} \times \frac{\Gamma_2(1 - q|\tau)\Gamma_2(\tau|\tau)}{\Gamma_2(1|\tau)\Gamma_2(\tau - q|\tau)}. \]  
It remains to observe that the functional equation of the double gamma function implies the identity
\[ \tau^{-\frac{q}{2}} \frac{\Gamma_2(1 - q|\tau)\Gamma_2(\tau|\tau)}{\Gamma_2(1|\tau)\Gamma_2(\tau - q|\tau)} = \frac{\Gamma(1 - q)}{\Gamma(1 - \frac{q}{2})}, \]  
which gives the result.

7 Conclusion and Open Questions

We introduced a new family \(\beta_{M,M-1}(a,b), M \in \mathbb{N}\), of Barnes beta distributions, i.e., distributions whose Mellin transform is defined to be an intertwining product of ratios of multiple gamma functions of Barnes, that are supported on \((0, \infty)\). We showed infinite divisibility and absolute continuity of \(\log \beta_{M,M-1}(a,b)\) by computing its Levy-Khinchine decomposition explicitly. We calculated the asymptotic expansion of the Mellin transform of \(\beta_{M,M-1}(a,b)\) and solved the Stieltjes moment problem for \(\beta_{M,M-1}(a,b)\). We also showed that the ratio \(\beta_{M,M-1}(a,b)\beta^{-1}_{M,M-1}(a,b)\) with different values of \(b_0\) possesses remarkable infinite factorizations and its Mellin transform reduces in a special case to intertwining products of ratios of multiple sine functions. For application, we noted that \(\beta_{2,1}(a,b)\) coincides with the distribution, which was first discovered in [26] and [27] and recently studied in [28] in the context of laws of certain functionals of stable processes, and the distribution \(\beta_{1,0}(a,b)\) appears as a building block of both Selberg and Morris probability distributions.

We reviewed the general theory of Barnes beta distributions to make our paper self-contained and then focused on the structure of the Selberg and Morris integral probability distributions. The construction of the latter is original to this paper. In both cases we emphasized the connection with Barnes beta distributions of types \((1,0)\) and \((2,2)\) and showed as a corollary that the Stieltjes moment problem for the negative moments of the Morris integral distribution, unlike that of the Selberg integral distribution, is determinate. We noted that the self-duality of the Mellin transform of the Selberg integral distribution, which was established originally in [18] and [19], follows from its involution invariance that we established in [39]. We showed that the self-duality of the Mellin transform of the Morris integral distribution, which was established originally in [19] in the special case of the Dyson integral distribution, extends to the general Morris integral distribution and also follows from its involution invariance, which we proved in this paper. For application, we examined the conjectures of [14], [18], and [19] about the maximum distribution of the gaussian free field on the interval and circle. We reviewed their calculations from the viewpoint of the gaussian multiplicative chaos theory and conjecturally expressed the law of the derivative martingale in both cases in terms of Barnes beta distributions of types \((1,0)\) and \((2,2)\). In particular, in the interval case, our probabilistic reformulation of their conjecture is new and, in the circle case, our conjecture itself is new and extends the original conjecture of [14] by adding a non-random logarithmic potential to the gaussian free field.
There are at least two additional areas of applications of Barnes beta distributions that are beyond the scope of this paper and left to future research. First, the Selberg integral distribution appears conjecturally as a mod-gaussian limit of the exponential functional of a large class of known statistics that converge to a gaussian process with $\mathcal{N}^{1/2}(\mathbb{R})$ limiting covariance, see [40]. Hence, the considerations of this paper can be naturally carried over to the maximum distribution of all such statistics. Second, we showed in [39] that a limit of Barnes beta distributions of type $\beta_{2M,3M}$ as $M \to \infty$ can be used to approximate the completed Riemann zeta function. It remains to see if this approximation is useful in verifying the conjectures of [16] about extremal values of the Riemann zeta function on the critical line, which are related to the Selberg and Morris integral distributions, i.e. if the aforementioned limit can be related to $\beta_{2,2}$ distributions.

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References

[1] E. Bacry, J. Delour, J.-F. Muzy: Multifractal random walk. Phys. Rev. E 64, 026103 (2001).
[2] E. Bacry, J.-F. Muzy: Log-infinitely divisible multifractal random walks. Comm. Math. Phys. 236, 449–475 (2003).
[3] E. W. Barnes: The genesis of the double gamma functions. Proc. London Math. Soc. s1-31, 358–381 (1899).
[4] E. W. Barnes: On the theory of the multiple gamma function. Trans. Camb. Philos. Soc. 19, 374–425 (1904).
[5] P. Biane, J. Pitman, M. Yor: Probability laws related to the Jacobi theta and Riemann zeta functions, and brownian excursions. Bulletin of the American Mathematical Society 38, 435–465 (2001).
[6] P. Bourgade and J. Kuan: Strong Szegő asymptotics and zeros of the zeta function. Comm. Pure Appl. Math 67, 1028–1044, http://arxiv.org/abs/1203.5328 (corrected version) (2013).
[7] X. Cao, A. Rosso, R. Santachiara: Extreme value statistics of 2D Gaussian free field: effect of finite domains. J. Phys. A: Math. Theor. 49, 02LT02 (2016).
[8] D. Carpentier and P. Le Doussal: Glass transition of a particle in a random potential, front selection in nonlinear renormalization group, and entropic phenomena in Liouville and sinh-Gordon models. Phys. Rev. E 63, 026110 (2001).
[9] C. Chamon, C. Mudry, X.-G. Wen: Localization in two dimensions, Gaussian field theories, and multifractality. Phys. Rev. Lett. 77, 4194 (1996).
[10] J. Ding, R. Roy, O. Zeitouni: Convergence of the centered maximum of log-correlated Gaussian fields, http://arxiv.org/abs/1503.04588 (2015).
[11] D. Dufresne: \textit{G} distributions and the beta-gamma algebra. \textit{Elect. J. Prob.} \textbf{15}, 2163–2199 (2010).

[12] B. Duplantier, S. Sheffield: Liouville quantum gravity and KPZ. \textit{Invent. Math.} \textbf{185}, 333-393 (2011).

[13] P. J. Forrester: Log-Gases and Random Matrices. \textit{Princeton University Press}, Princeton (2010).

[14] Y. V. Fyodorov and J. P. Bouchaud: Freezing and extreme-value statistics in a random energy model with logarithmically correlated potential. \textit{J. Phys. A, Math Theor.} \textbf{41}, 372001 (2008).

[15] Y. V. Fyodorov, O. Giraud: High values of disorder-generated multifractals and logarithmically correlated processes. \textit{Chaos, Solitons & Fractals}, \textbf{74}, 15–26 (2015).

[16] Y. V. Fyodorov and J. P. Keating: Freezing transitions and extreme values: random matrix theory, $\zeta(1/2+it)$, and disordered landscapes. \textit{Philos. Trans. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci.} \textbf{372}, 20120503 (2014).

[17] Y. V. Fyodorov, B. A. Khoruzhenko, and N. J. Simm: Fractional Brownian motion with Hurst index $H = 0$ and the Gaussian Unitary Ensemble. \url{http://arxiv.org/abs/1312.0212} (2015).

[18] Y. V. Fyodorov, P. Le Doussal: Moments of the position of the maximum for GUE characteristic polynomials and for log-correlated Gaussian processes. \textit{J. Stat. Phys.} DOI 10.1007/s10955-016-1536-6, 1–51 (2016).

[19] Y. V. Fyodorov, P. Le Doussal, A. Rosso: Statistical mechanics of logarithmic REM: duality, freezing and extreme value statistics of $1/f$ noises generated by gaussian free fields. \textit{J. Stat. Mech. Theory Exp.}, P10005 (2009).

[20] Y. V. Fyodorov, P. Le Doussal, A. Rosso: Counting function fluctuations and extreme value threshold in multifractal patterns: the case study of an ideal $1/f$ noise. \textit{J. Stat. Phys.} \textbf{149}, 898–920 (2012).

[21] Y. V. Fyodorov, N. J. Simm: On the distribution of maximum value of the characteristic polynomial of GUE random matrices, \url{http://arxiv.org/abs/1503.07110} (2015).

[22] C. P. Hughes, J. P. Keating, and N. O’Connell: On the characteristic polynomial of a random unitary matrix. \textit{Commun. Math. Physics.} \textbf{220}, 429–451 (2001).

[23] J. Jacod, E. Kowalski, A. Nikeghbali: Mod-gaussian convergence: new limit theorems in probability and number theory. \textit{Forum Math.} \textbf{23}, 835–873 (2011).

[24] J.-P. Kahane: Positive martingales and random measures. \textit{Chinese Ann. Math. Ser. B} \textbf{8}, 1–12 (1987).

[25] N. Kurokawa, S. Koyama: Multiple sine functions. \textit{Forum Math.} \textbf{15}, 839–876 (2003).

[26] A. Kuznetsov: On extrema of stable processes. \textit{Ann. Probab.} \textbf{39}, 1027–1060 (2011).

[27] A. Kuznetsov, J. C. Pardo: Fluctuations of stable processes and exponential functionals of hypergeometric Lévy processes. \textit{Acta Appl. Math.} \textbf{123}, 113–139 (2013).
26

[28] J. Letemplier, T. Simon: On the law of homogeneous stable functionals, http://arxiv.org/abs/1510.07441 (2015).

[29] G. D. Lin, J. Stoyanov: Moment determinacy of powers and products of nonnegative random variables. J. Theor. Probab. 28, 1337–1353 (2015).

[30] T. Madaule: Maximum of a log-correlated Gaussian field, http://arxiv.org/abs/1307.1365 (2014).

[31] T. Madaule, R. Rhodes, V. Vargas: Glassy phase and freezing of log-correlated Gaussian potentials. Ann. Appl. Probab. 26, 643-690 (2016).

[32] B. B. Mandelbrot: Possible refinement of the log-normal hypothesis concerning the distribution of energy dissipation in intermittent turbulence, in Statistical Models and Turbulence, M. Rosenblatt and C. Van Atta, eds., Lecture Notes in Physics 12, Springer, New York, p. 333 (1972).

[33] B. B. Mandelbrot: Limit lognormal multifractal measures, in Frontiers of Physics: Landau Memorial Conference, E. A. Gotsman et al, eds., Pergamon, New York, p. 309 (1990).

[34] J.-F. Muzy, E. Bacry: Multifractal stationary random measures and multifractal random walks with log-infinitely divisible scaling laws. Phys. Rev. E 66, 056121 (2002).

[35] A. Nikeghbali and M. Yor: The Barnes G function and its relations with sums and products of generalized gamma convolutions variables. Elect. Comm. in Prob. 14, 396–411 (2009).

[36] D. Ostrovsky: Mellin transform of the limit lognormal distribution. Comm. Math. Phys. 288, 287–310 (2009).

[37] D. Ostrovsky: Selberg integral as a meromorphic function. Int. Math. Res. Not. IMRN, 17, 3988–4028 (2013).

[38] D. Ostrovsky: Theory of Barnes beta distributions. Elect. Comm. in Prob. 18, no. 59, 1–16 (2013).

[39] D. Ostrovsky: On Barnes beta distributions, Selberg integral and Riemann xi. Forum Math., DOI: 10.1515/forum-2013-0149 (2014).

[40] D. Ostrovsky: On Riemann zeroes, lognormal multiplicative chaos, and Selberg integral. Nonlinearity 29, 426–464 (2016).

[41] B. S. Rajput, J. Rosinski: Spectral representations of infinitely divisible processes. Probab. Theory Relat. Fields 82, 451–487 (1989).

[42] R. Rhodes and V. Vargas: Gaussian multiplicative chaos and applications: an overview. Probability Surveys 11, 315–392 (2014).

[43] B. Rodgers: A central limit theorem for the zeroes of the zeta function. Int. J. Number Theory, 10, 483–511 (2014).

[44] S. N. M. Ruijsenaars: On Barnes’ multiple zeta and gamma functions. Adv. Math. 156, 107–132 (2000).
[45] T. Shintani: A proof of the classical Kronecker limit formula. *Tokyo J. Math.* **3**, 191–199 (1980).

[46] A. Soshnikov: The central limit theorem for local linear statistics in classical compact groups and related combinatorial identities. *Ann. Prob.*, **28**, 1353–1370 (2000).

[47] F. W. Steutel and K. van Harn: Infinite Divisibility of Probability Distributions on the Real Line. *Marcel Dekker*, New York (2004).

[48] E. Subag, O. Zeitouni: Freezing and Decorated Poisson Point Processes. *Commun. Math. Phys.* **337**, 55–92 (2015).

[49] C. Webb: The characteristic polynomial of a random unitary matrix and Gaussian multiplicative chaos - The $L^2$-phase. *Elect. J. Prob.* **20**, no. 104, 1–21 (2015).