Heat Transfer in Cadmium Telluride-Water Nanofluid over a Vertical Cone under the Effects of Magnetic Field inside Porous Medium

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Abstract: The present research provides a numerical investigation of two dimensional nanofluid flow over an inverted cone inside a porous medium. The model is developed to incorporate non-spherical shapes of \( \text{CdTe} \)-nanoparticles in water based fluid. Simultaneous effects of pertinent parameters like volume fraction, Reynold number, Hartmann number, porosity, Grashof number, radiation parameter and Peclet number on temperature distribution and velocity profile are studied and illustrated graphically. In addition, the corresponding computational results of Nusselt number and skin friction for regulating parameters are also presented in graphs and tables. The highest Nusselt number is observed for blade-shaped \( \text{CdTe} \) particles. Furthermore, the thermal conductivity and viscosity are also calculated for non-spherical shapes of \( \text{CdTe} \) nanoparticles. The result showed that the thermal conductivity of nanofluid with blade-shaped particles is 0.94% and 1.93% greater than platelet and brick type particles. The computational results for the special case are validated by comparisons with the presented results in previous studies and the results are in perfect agreement.

Keywords: magnetohydrodynamics; nanofluid; vertical cone; finite difference method

1. Introduction

A stable crystalline/powder cadmium telluride (\( \text{CdTe} \)) is formed with cadmium (metallic) and telluride (crystalline). Margottet was the first who introduced \( \text{CdTe} \) while reacting metals with telluride at red heat [1]. The bonding structure of \( \text{CdTe} \) is shown in Figure 1. \( \text{CdTe} \) plays important role in several industrial applications. It has magnificent properties like high optical absorption, the ability to absorb sunlight at ideal length, band gap (1.5 ev), abundant. \( \text{CdTe} \) thin films are commercially used in solar cells due to its magnificent properties. \( \text{CdTe} \) thin film (2 \( \mu \)m) has capability capture 100% solar radiation [2]. Owing to that fact, \( \text{CdTe} \)-solar cells are the leading candidate in photovoltaic development. \( \text{CdTe} \)-solar cells have been discussed by many researchers. Furthermore, their thermal and electrical properties have also attracted the attention of researchers [3–7]. However, researchers did not pay attention to improving the thermal performance of weak convectional fluid using \( \text{CdTe} \) nanoparticles. To the best of authors’ knowledge, for the first time, the heat transfer intensification in ethylene glycol by suspension of \( \text{CdTe} \) nanoparticles is worked out by Hanif et al. [8].
Figure 1. Bonding structure of CdTe.

In thermal systems, the properties of convectional heat transfer fluid are neither sufficient nor effective due to their low thermal conductivity. Therefore, researchers are more interested in nanofluid compared with regular convectional fluid since they are capable to improve the efficiency of thermal system due to increased thermal conductivity. A nanofluid is a colloidal mixture of regular fluid and nanoparticles. It has been shown that certain combinations of nano-sized particles and fluids possessed comparatively higher thermal rates than the regular fluid. The potential uses of nanofluid are typically in the heat transfer region, where a small concentration of nanoparticles often intensifies the heat transfer rates, for instance, see References [9–14].

Nanofluids have unique properties like viscosity and thermal conductivity and these properties vary by changing base fluid, size, shape and volume fraction of nanoparticles. Over the past few years, several studies have been done using different shapes and sizes of nanoparticles for the purpose of better thermal performance in thermal systems. Khan [15] discussed different shapes of MoS$_2$-water based nanofluid in the presence of a magnetic field. It was concluded that blade-shaped particles gained 8.1%, 4.8%, 4%, 3.4% and 0.2% greater heat transfer rates than spherical, brick, cylinder and platelet shaped of MoS$_2$ − H$_2$O nanofluid. Hassan et al. [16] compared the heat transfer rates of three different shapes (prolate, oblate and sphere) of Fe$_3$O$_4$ nanoparticles in water based nanofluid over a rotating disc. They found that maximum temperature rates are attained by prolate shaped nanoparticles. Moreover, the temperature rates are higher when a magnetic field is considered. Magnetohydrodynamics (MHD) force convection in water based nanofluid, considering different shapes of CuO particles, is studied numerically by sheikholeslami et al. [17]. The results showed that platelet shape particles gained maximum heat transfer rates. Also, the temperature gradient decreases the function of the Hartmann number. Sheikholeslami et al. [18] studied the shape variation in Al$_2$O$_3$ nanofluid in a permeable enclosure with MHD effects using the Darcy model. Hassan et al. [19] discussed the shape effects of iron nanoparticles in convectional ferrofluid flow over a rotating disk in the presence of an oscillating magnetic field. Hamid et al. [20] explored the shape effects of MoS$_2$ nanoparticles on rotating fluid flow through a stretching elastic surface. Sheikholeslami et al. [21] investigated the escalation of heat transfer in ferrofluid inside a porous enclosure under the shape effects of Fe$_3$O$_4$ particles. Akbar et al. [22] studied the shape effects of nano-scale shape particles on the peristaltic flow of water based nanofluid in combination with magnetic field. Shi et al. [23] analyzed heat transfer combined with entropy generation in nanofluid flow through a channel using different shapes of CuO-nanoparticles. Kumar et al. [24] investigated shape effects of Cu nanoparticles in water based nanofluid in a channel. They found an augmentation in Nusselt number due to the increment in thermal radiation and Rayleigh number, whereas a reverse trend was found by increasing the Hartmann number. Some interesting experimental and theoretical work on nanoparticle shape effects can be found in References [25–31].

In several real-life situations, fluid flow over a cone is in high demand, for example in healthcare systems, aeronautical engineering, energy storage, astrophysics, space technology, and so forth. Also, in chemical industry, canonical surfaces have been used in filtering, pumping, drilling and degassing machines, and so forth. It is therefore important not only to research these parameters, but also to have an idea of the behavior of liquid fluid flow when it flows along this shape in particular. It is
also important to understand and forecast the flow nature and the involvement of nanoparticles in applications for heat transfer to be used in industrial chemical processes.

A literature survey reveals that the frequently used nanoparticles in nano thermal sciences are Al\(_2\)O\(_3\), Cu and carbon nanotubes (CNTs) and so forth. Also, spherical nanoparticles are the most common particles considered by researchers but, when it comes to their application and significance, they are limited. Owing to the fact, this study incorporates the non-spherical CdTe nanoparticles. More precisely, blade, brick and platelet shapes are taken into account. We believe a numerical study on different particle shapes of CdTe contained in water based fluid has not yet been reported. Therefore, this is the first attempt to analyse heat transfer rates in water based nanofluid under the shape effects of CdTe nanoparticles. Furthermore, radiation effects on natural convection flow over a vertical cone in the presence of a magnetic field are also discussed. The thermo-physical properties of base fluid water and CdTe nanoparticles are tabulated in Table 1. The effects of regulating parameters \(\phi\), \(K\), \(Ha\), \(Gr\) and \(Rd\) on temperature distribution, velocity profile are illustrated graphically and discussed in detail. Furthermore, the numerical results for Nusselt number \(Nu\) and skin friction \(\tau\) under the influence of active parameters are also tabulated. For validation purposes, the computed results for some special cases are compared with the results of previously published work.

### Table 1. Thermal and mechanical properties of water [15] and CdTe nanoparticles.

| Materials          | Water | CdTe | Reference |
|--------------------|-------|------|-----------|
| \(\rho\) (kg/m\(^3\)) | 997.1 | 5855 | [32]      |
| \(C_p\) (J/kgK)       | 4179  | 209  | [33]      |
| \(k\) (W/mK)          | 0.613 | 7.5  | [34]      |
| \(\beta\) (K\(^{-1}\)) | \(21 \times 10^{-5}\) | \(0.5 \times 10^{-5}\) | [34] |
| \(\sigma = 1/\text{resistivity (Sm}^{-1}\right)\) | \(0.55 \times 10^{-5}\) | \(0.2 \times 10^{-7}\) | [35] |

### 2. Mathematical Formulation

A free convection in water based CdTe nanofluid over an inverted cone inside a porous medium is considered. In addition, three different shapes of CdTe nanoparticles, namely, blade, brick and platelet are taken into account. The surface of the cone is taken as \(x - axis\) and normal to the cone is referred to \(y - axis\).

A uniform magnetic field of strength \(B_0\) is considered normal to \(x - axis\), see Figure 2. The nanofluid is supposed to be electrically conducting due to an applied magnetic field. The induced magnetic field can be neglected since the Reynold number is sufficiently small. Initially (\(\bar{t} \leq 0\)), the temperature distribution of the cone surface is considered the same as the surrounding fluid, that is, \(T_\infty\). Thereafter (\(\bar{t} > 0\)), temperature \(T\) of the nanofluid at the surface of cone increased to \(T_w = T_\infty + cx^n\). The 2D unsteady, unidirectional and incompressible nanofluid together with Boussinesq approximation can be written as

\[
\nabla \cdot (rv) = 0,
\]

\[
\rho_{nf} \left( \frac{\partial \bar{u}}{\partial \bar{t}} + (v \cdot \nabla) \bar{u} \right) = \mu_{nf} \frac{\partial^2 \bar{u}}{\partial \bar{y}^2} - \sigma_{nf} B_0^2 \bar{u} - \frac{\mu_{nf}}{k_0} \bar{u} + g(\rho \beta)_{nf} (\bar{T} - T_\infty) \cos \phi,
\]

\[
(\rho C_p)_{nf} \left( \frac{\partial \bar{T}}{\partial \bar{t}} + (v \cdot \nabla) \bar{T} \right) = k_{nf} \frac{\partial^2 \bar{T}}{\partial \bar{y}^2} - \frac{\partial q_r}{\partial \bar{y}}.
\]
where \( \mathbf{v} = (\bar{u}(\bar{x}, \bar{y}, t), \bar{v}(\bar{x}, \bar{y}, t)) \) and \( \bar{T} = \bar{T}(\bar{x}, \bar{y}, t) \) represent 2-dimensional velocity in \((\bar{x}, \bar{y})\) – plane and temperature receptively. The mathematical expressions for nanofluid properties are tabulated in Table 2. The appropriate initial boundary conditions are:

\[
\begin{align*}
\bar{t} \leq 0 : & \quad \bar{u} = 0, \quad \bar{v} = 0, \quad \bar{T} = T_\infty, \\
\bar{t} > 0 : & \quad \bar{u} = 0, \quad \bar{v} = 0, \quad \bar{T} = T_\infty + c\bar{x}^n \quad \text{at} \quad \bar{y} = 0, \\
& \quad \bar{u} = 0, \quad \bar{T} = T_\infty \quad \text{at} \quad \bar{x} = 0, \\
& \quad \bar{u} \to 0, \quad \bar{T} \to T_\infty \quad \text{as} \quad \bar{y} \to \infty.
\end{align*}
\]

\( \text{(4)} \)

Figure 2. Problem schematics and geometrical coordinates.

Table 2. Mathematical expressions for nanofluid properties.

| Properties          | Nanofluid \( \rho_{nf} = (1 - \varphi)\rho_f + \varphi\rho_s \) | References |
|---------------------|---------------------------------------------------------------|------------|
| Density             | \( \mu_{nf} = \mu_f(1 + a_0\varphi + b_0\varphi^2) \) \quad \text{where} \quad a_0 \quad \text{and} \quad b_0 \quad \text{are shape constants given in Table 3.} | \[15\]      |
| Viscosity           | \( \rho = \rho_f(1 - \varphi)\rho_f + \varphi\rho_s \)      | \[37\]      |
| Thermal expansion   | \( (\rho\beta)_{nf} = (1 - \varphi)(\rho\beta)_f + \varphi(\rho\beta)_s \) | \[37\]      |
| Heat capacitance    | \( (\rho C_p)_{nf} = (1 - \varphi)(\rho C_p)_f + \varphi(\rho C_p)_s \) | \[37\]      |
| Thermal conductivity| \( \frac{k_{nf}}{k_f} = \frac{(k_s + (m - 1)k_f)}{(k_s + (m - 1)k_f) - \varphi(k_s - k_f)} \) \quad \text{with} \quad m = 3/\psi_0 \quad \text{where} \quad \psi_0 \quad \text{is the sphericity and given in Table 3.} | \[24\]      |
| Electrical conductivity | \( \frac{\sigma_{nf}}{\sigma_f} = 1 + \frac{3\varphi}{\left(\frac{\sigma_s}{\sigma_f} + 2\right) - \varphi\left(\frac{\sigma_s}{\sigma_f} - 1\right)} \) | \[37\]      |
Table 3. The shape constants \( a_0, b_0 \) and sphericity \( \psi_0 \) of nanoparticles [15].

| Particle Category | Shape Constants | Sphericity | Shape |
|-------------------|-----------------|------------|-------|
| Blade             | 14.6, 123.3     | 0.36       |       |
| Brick             | 1.9, 471.4      | 0.81       |       |
| Platelet          | 37.1, 612.6     | 0.52       |       |

Rosseland approximation for radiative heat flux is (see Reference [38]):

\[
q_r = \frac{4\epsilon_b}{3k_b} \frac{\partial T^4}{\partial y^*},
\]

where \( \epsilon_b \) and \( k_b \) denote Stefan–Boltzman and the absorption coefficient, respectively. Suppose the difference \( \bar{T} - T_\infty \) inside flow is sufficiently small and \( T^4 \) can be expanded by about \( T_\infty \) by using the Taylor series. Hence, the approximation of \( T^4 \) by neglecting higher orders:

\[
T^4 \approx T_\infty^4 + 4T_\infty^3(T - T_\infty).
\]

Thereafter, using Equations (5) and (6) into Equation (3) leads to

\[
(\rho C_p)_n f \left( \frac{\partial T}{\partial t} + (v \cdot \nabla) \bar{T} \right) = \left( k_n f + \frac{16\epsilon_b}{3k_b} T_\infty^3 \right) \frac{\partial^2 T}{\partial y^2}.
\]

Invoking the following appropriate non-dimensional quantities

\[
x = \frac{\bar{x}}{L}, \quad y = \frac{\bar{y}}{L}, \quad r = \frac{\bar{r}}{L}, \quad t = \frac{\bar{t} u_0}{L},
\]

\[
u = \frac{\bar{u}}{u_0}, \quad \bar{v} = \frac{\bar{v}}{u_0}, \quad \bar{T} = \frac{(\bar{T} - T_\infty)}{(T_w - T_\infty)}, \quad \bar{\tau}_x = \frac{\tau_5 L}{\mu_f u_0},
\]

in Equations (1), (2) and (7) together with nanofluid properties defined in Table 2. We have,

\[
\frac{\partial}{\partial x} (ru) + \frac{\partial}{\partial y} (rv) = 0,
\]

\[
\phi_1 Re \left[ \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right] = \phi_2 \frac{\partial^2 u}{\partial y^2} - \phi_2 \frac{1}{K} u - \phi_3 H u^2 u + \phi_4 Gr T \cos \phi,
\]

\[
\phi_5 Re \left[ \frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right] = (\phi_6 + Rd) \frac{\partial^2 T}{\partial y^2},
\]
where

\[ \phi_1 = (1 - \varphi) + \varphi \frac{\rho_f}{\rho_f}^s, \quad \phi_2 = 1 + a_0 \varphi + b_0 \varphi^2, \quad \phi_3 = \frac{\sigma_{nf}}{\sigma_f}, \]

\[ \phi_4 = (1 - \varphi) + \varphi \frac{(\rho \beta)_s}{(\rho \beta)_f}^s, \quad \phi_5 = (1 - \varphi) + \varphi \frac{(\rho C_P)_s}{(\rho C_P)_f}^s, \quad \phi_6 = \frac{k_{nf}}{k_f}, \]

\[ Re = \frac{u_0 L}{v_f}, \quad \frac{1}{K} = \frac{L^2}{k_f}, \quad Ha = B_0 L \sqrt{\frac{\sigma_f}{\mu_f}}, \]

\[ Gr = \frac{g \beta_f (T_w - T_\infty) L^2}{v f u_0}, \quad Pe = \frac{u_0 L (\rho C_P)_f}{k_f}, \quad Rd = \frac{16 c_0 T_\infty^3}{3 k_b k_f}. \]

Here, \( \phi_1 - \phi_6 \) refer to nanofluid constants. \( Re, K, Ha, Gr, Rd, \) and \( Pe \) represent the Reynold’s number, permeability parameter, Hartmann number, Grashof number, radiation parameter, and Peclet number respectively. The dimensionless form of Equation (4) is:

\[
\begin{align*}
    t \leq 0 : & \quad u = 0, \quad v = 0, \quad T = 0, \\
    t > 0 : & \quad u = 0, \quad v = 0, \quad T = x^n \quad \text{at} \quad y = 0, \\
    & \quad u = 0, \quad T = 0 \quad \text{at} \quad x = 0, \\
    & \quad u \to 0, \quad T \to 0 \quad \text{as} \quad y \to \infty.
\end{align*}
\]

3. Nusselt Number and Skin Friction

Inducing non-dimensional parameters Equation (8) in Local Nusselt number \( \overline{Nu}_x \) and skin friction coefficient \( \overline{\tau}_x \) defined below,

\[ \overline{Nu}_x = \frac{1}{k_f (T_w - T_\infty)} \left[ -k_{nf} x \left( \frac{\partial T}{\partial y} \right)_{y=0} + (q_f)_{y=0} \right], \quad \overline{\tau}_x = \mu_{nf} \left( \frac{\partial u}{\partial y} \right)_{y=0}. \]

We have,

\[ \overline{Nu}_x = -x (\phi_6 + Rd) \left( \frac{\partial T}{\partial y} \right)_{y=0}, \quad \overline{\tau}_x = \phi_2 \left( \frac{\partial u}{\partial y} \right)_{y=0}. \]

4. Numerical Analysis

The resulting system of non-linear coupled Equations (9)–(11) along with corresponding conditions Equation (13) is discretized by a finite difference scheme, namely, the Crank Nicolson method. It is an implicit, accurate, fast convergent, and an unconditionally stable method. This scheme approximates the equations at an average of current and forward time level, that is, \( n^{th} \) and \( (n+1)^{th} \) time level.

The rectangular area of integration has limits \( x = 1 \) and \( y = 15 \). Here \( y \) corresponds to \( y = \infty \) which positions quite away from exterior of the momentum and thermal boundary layers. The time step is taken as small as \( \Delta t = 0.01 \) combined with the mesh size \( \Delta x = 0.05, \Delta y = 0.05 \) in \( x \) and \( y \) direction, respectively. For the convergent solution, the numerical iterations are repeated for several time. The solution is supposed to be converged when an absolute error approaches \( 1e^{-5} \) for all grid nodes.

5. Graphical Results and Discussion

The different shapes of \( CdTe \) particles in the MHD flow of cadmium telluride nanofluid past a vertical cone with thermal radiation effects are considered. The Hamilton and Crosser model
[39] is used for the thermal conductivity of different shapes of CdTe nanoparticles. For validation purposes, some limited cases of the current study are compared with the numerical results provided in References [40,41], as shown in Table 4 and the results are in excellent agreement.

Variational effects of of regulating parameters particle volume fraction ($\phi = 0, 0.01, 0.02, 0.03$), porosity parameter ($K = 0.5, 1.5, 2.5, 3.5$), Hartmann number ($M = 0, 2, 4, 6$), thermal Grashof number ($Gr = 0.1, 0.3, 0.5, 0.7$), and thermal radiation ($Rd = 0, 1, 2, 3$) on the temperature and velocity profiles are plotted graphically. Additionally, the influence of regulating parameters on physical quantities $Nu_x$ and $\tau_x$ are also calculated and tabulated in Tables 5–7. The results are reasonable and more close to physical expectations.

Figures 3–12 have plotted for different shapes of CdTe nanoparticles under the influence of emerging parameters ($\phi, K, Ha, Gr$ and $Rd$) at fixed parameters $\phi = \pi/4$, $Re = 1$, $Pe = 0.3$ and $n = 0.5$. Figures 3 and 4 depict the change in temperature and velocity field for different shapes of CdTe nanoparticles under the effect of particle volume fraction $\phi$. It is perceived that with an augment in $\phi$ the temperature distribution increases. Whereas the velocity field showed an opposite behaviour. The maximum increase in temperature and minimum decrease in velocity are observed for platelet-shaped nanoparticles followed by blade and brick-shaped nanoparticles.

| $Pe$ | $Nu_x$ | $\tau_x$ |
|------|--------|----------|
|      | Present | [40] [41] |
| 0.01 | 0.0749  | 0.0751   | 0.0884  | 1.3551 | 1.3549 | 1.3158 |
| 0.1  | 0.2116  | 0.2116   | 0.2109  | 1.0960 | 1.0962 | 1.0985 |
| 1    | 0.5109  | 0.5111   | 0.5119  | 0.7699 | 0.7697 | 0.7696 |
| 10   | 1.0339  | 1.0342   | 1.0337  | 0.4877 | 0.4877 | 0.4861 |
| 100  | 1.9226  | 1.9230   | 1.9235  | 0.2896 | 0.2895 | 0.2882 |

Table 4. Comparison of steady state local Nusselt number and skin friction at $x = 1$ for different values of $Pe$ when $\phi = 1/K = Ha = Rd = n = 0$, and $Re = Gr = 1$.

Figure 3. Influence of $\phi$ on temperature distribution when $K = 0.5$, $Ha = 2$, $Gr = 0.3$, and $Rd = 3$. 
Figure 4. Influence of $\varphi$ on velocity profile when $K = 0.5, Ha = 2, Gr = 0.3$, and $Rd = 3$.

Figure 5. Influence of $K$ on temperature distribution when $\varphi = 0.01, Ha = 2, Gr = 0.3$, and $Rd = 3$.

Figure 6. Influence of $K$ on velocity profile when $\varphi = 0.01, Ha = 2, Gr = 0.3$, and $Rd = 3$. 
Figure 7. Influence of $Ha$ on temperature distribution when $\varphi = 0.01$, $K = 0.5$, $Gr = 0.3$, and $Rd = 3$.

Figure 8. Influence of $Ha$ on velocity profile when $\varphi = 0.01$, $K = 0.5$, $Gr = 0.3$, and $Rd = 3$.

Figure 9. Influence of $Gr$ on temperature distribution when $\varphi = 0.01$, $K = 0.5$, $Ha = 2$, and $Rd = 3$. 
Figure 10. Influence of $Gr$ on velocity profile when $\varphi = 0.01$, $K = 0.5$, $Ha = 2$, and $Rd = 3$.

Figure 11. Influence of $Rd$ on temperature distribution when $\varphi = 0.01$, $K = 0.5$, $Ha = 2$, and $Gr = 0.3$.

Figure 12. Influence of $Rd$ on velocity profile when $\varphi = 0.01$, $K = 0.5$, $Ha = 2$, and $Gr = 0.3$. 
Figures 5 and 6 designate the effects of porosity parameter $K$ on temperature distribution and velocity profile, respectively. It is noticed that temperature distribution and velocity field decreases and increases respectively with an increase in porosity parameter. The brick-shaped particles possess lowest temperature and highest velocity rates followed by blade-shaped and platelet-shaped particles.

The behavior of Hartmann number $Ha$ in the presence of three different shapes of CdTe nanoparticles on the temperature distribution and velocity profile of nanofluid are illustrated in Figures 7 and 8, respectively. The temperature distribution rises with an augment in the value of $Ha$ whereas the velocity profile dropped off for higher values of $Ha$.

In Figure 9, the variation of the thermal Grashof number $Gr$ is examined on the temperature distribution of water based CdTe-nanofluid. It is noticed that the temperature distributions decreases by increasing the value of $Gr$ and a maximum dropped down in the temperature graph is attained by brick-shaped particles. The velocity change due to $Gr$ is plotted in Figure 10. For higher values of $Gr$, a significant increase is observed in fluid flow. It is reasonable in the sense that maximum Grashof number is a source to upgrade the buoyancy force which buoyed up the nanofluid flow rates as a result the velocity field increases.

Figures 11 and 12 are plotted to examine the temperature distribution and velocity profile of CdTe nanofluid against the variation in the radiation parameter $Rd$. From the plots, it is evident that the temperature and velocity profiles increase when $Rd$ is increased. This behavior is closed to physical expectations as the energy production causes elevation of the temperature and fluid flow.

Figures 13 and 14 display the variation in Nusselt number $Nu_x$ and skin friction $\tau_x$ due to volume fraction $\varphi$ of CdTe nanoparticles. It is noticed that both $Nu_x$ and $\tau_x$ increase when $\varphi$ increases. It is important to bear in mind that blade-shaped particles have maximum sphericity compared to others. For the reason, the blade-shaped particles gained maximum values for $Nu_x$. Platelet-shaped particles attained minimum Nusselt number for ($0 < \varphi < 0.035$) than brick-shaped particles. On the other hand, maximum values for $\tau_x$ are obtained by platelet-shaped particles compared with blade and brick-shaped particles. The blade-shaped particles showed the highest rates of skin friction than brick-shaped particles when ($\varphi < 0.035$).

![Figure 13. Influence of $\varphi$ on Nusselt number when $K = 0.5$, $Ha = 2$, $Gr = 0.3$ and $Rd = 3$.](image-url)
In Figure 14, the thermal conductivity is plotted against CdTe particle volume fraction. It is evident from the plot that the thermal conductivity for blade-shaped particle increases maximum with respect to volume fraction than other particles. The result is sensible in a way that the blade type particle has the highest sphericity compared with platelet and brick particles. On the other hand, the values of shape constant $a_0$ and $b_0$ for platelet particles are higher than that of blade and brick. Owing to the fact, platelet type particles possessed maximum viscosity, as depicted in Figure 16.

Figure 14. Influence of $\varphi$ on Skin friction when $K = 0.5, Ha = 2, Gr = 0.3$ and $Rd = 3$.

Figure 15. Dependence of thermal conductivity on volume fraction of CdTe nanoparticles.
Figure 16. Dependence of viscosity on volume fraction of CdTe nanoparticles.

Tables 5–7 are tabulated to examine the effects of emerging parameters (K, Ha, and Rd) on Nusselt number \( Nu_x \) and skin friction \( \tau_x \). The effects of porosity parameter \( K \) on \( Nu_x \) and \( \tau_x \) are tabulated in Table 5. It is found that the maximum Nusselt number is attained by the nanofluid in presence of blade-shaped particles compared with brick and platelet particles. However, the maximum skin friction for maximum value of \( K \) is attained by platelet particles followed by blade and brick type particles.

**Table 5.** Nusselt number \((-Nu_x)\) and skin friction \((\tau_x)\) for different values of \( K \) at \( x = 1 \) when \( \phi = 0.01, Ha = 2, Gr = 0.3, \) and \( Rd = 3 \).

| \( K \) | Local Nusselt Number | Local Skin Friction |
|-------|----------------------|--------------------|
|       | Blade | Brick | Platelet | Blade | Brick | Platelet |
| 0.5   | 0.3159 | 0.3158 | 0.3118 | 0.0874 | 0.0851 | 0.0930 |
| 1.5   | 0.3271 | 0.3267 | 0.3242 | 0.1000 | 0.0967 | 0.1087 |
| 2.5   | 0.3305 | 0.3299 | 0.3276 | 0.1032 | 0.0996 | 0.1130 |
| 3.5   | 0.3317 | 0.3311 | 0.3295 | 0.1047 | 0.1010 | 0.1149 |

**Table 6.** Nusselt number \((-Nu_x)\) and skin friction \((\tau_x)\) for different values of \( Ha \) at \( x = 1 \) when \( \phi = 0.01, K = 0.5, Gr = 0.3, \) and \( Rd = 3 \).

| \( Ha \) | Local Nusselt Number | Local Skin Friction |
|--------|----------------------|--------------------|
|        | Blade | Brick | Platelet | Blade | Brick | Platelet |
| 0      | 0.03691 | 0.3743 | 0.3523 | 0.1388 | 0.1386 | 0.1394 |
| 2      | 0.3159 | 0.3158 | 0.3118 | 0.0874 | 0.0851 | 0.0930 |
| 4      | 0.2908 | 0.2896 | 0.2896 | 0.0522 | 0.0503 | 0.0571 |
| 6      | 0.2834 | 0.2821 | 0.2826 | 0.0362 | 0.0348 | 0.0400 |

Table 6 exhibits the impact of Hartmann number \( Ha \) on \( Nu_x \) and \( \tau_x \). A reduction in the magnitude of Nusselt number and the skin friction of CdTe nanofluid are observed for higher values of \( Ha \). It is also noticed that the nanofluid with brick-shaped particles showed the highest Nusselt number in the absence of a magnetic field, otherwise the highest values for the Nusselt number were obtained by
blade-shaped particles. Moreover, CdTe nanofluid possesses larger values of Nusselt number with brick-shaped nanoparticles as compared to platelet-shaped particles when $Ha < 4$ and an opposite relation is observed when $Ha > 4$. The skin friction also shows decreasing effects by intensifying the values of the Hartmann number and a minimum decrease is obtained by nanofluid suspended with brick-shaped nanoparticles.

Table 7. Nusselt number ($-Nu_x$) and skin friction ($\tau_x$) for different values of $Rd$ at $x = 1$ when $\varphi = 0.01$, $K = 0.5$, $Ha = 2$, and $Gr = 0.3$.

| $Rd$ | Local Nusselt Number | Local Skin Friction |
|------|----------------------|---------------------|
|      | Blade                | Brick               | Platelet            | Blade | Brick | Platelet |
| 0    | 0.1144               | 0.1140              | 0.1108              | 0.0862 | 0.0840 | 0.0917   |
| 1    | 0.1829               | 0.1828              | 0.1790              | 0.0869 | 0.0847 | 0.0925   |
| 2    | 0.2496               | 0.2495              | 0.2456              | 0.0872 | 0.0850 | 0.0928   |
| 3    | 0.3159               | 0.3158              | 0.3118              | 0.0874 | 0.0851 | 0.0930   |

Table 7 reflects the impact of radiation parameter on $Nu_x$ and $\tau_x$ and the result shows that both $Nu_x$ and $\tau_x$ are increasing functions of the radiation parameter. It is also observed that the blade and brick shaped nanoparticles showed almost the same increasing rates for Nusselt number in the presence of thermal radiation. Furthermore, minimum Nusselt number $Nu_x$ and maximum skin friction values $\tau_x$ are attained by platelet-shaped particles.

It is important to mention that different shapes of nano-sized particles only affect the thermal conductivity and viscosity of nanofluid. However, particle volume fraction not only enhances the thermal conductivity and viscosity of nanofluid, but also improves the values of other thermo-physical properties of fluid. Table 8 is drawn to display the numeric values of thermal conductivity $k_{nf}$ and viscosity $\mu_{nf}$ of CdTe nanofluid using models from References [39,42]. From the table, it is clearly seen that CdTe nanofluid with blade-shaped particles owned maximum thermal conductivity followed by platelet and brick type particles. Whereas the maximum viscosity of CdTe nanofluid is obtained by platelet-shaped nanoparticles. In addition, the viscosity of blade-shaped particle is greater than that of brick-shaped particles for $(0 < \varphi \leq 0.035)$. It is remarkable that the sequence of viscosity is followed by the velocity of different shapes of CdTe nanoparticles, the platelet-shaped particles having the highest viscosity gives a lower velocity profile. But the temperature profiles showed a different sequence than the thermal conductivity of water based CdTe nanofluid. A maximum thermal conductivity is attained in nanofluid containing blade-shaped CdTe nanoparticles followed by platelet and brick type particles. Whereas the blade-shaped and platelet-shaped particles change the sequence in the temperature profile. It is worth mentioning that the behavior of temperature and velocity are not only affected by thermal conductivity and viscosity, but also by other thermo-physical properties. However, temperature distribution increases with respect to the volume fraction of all active-shapes of CdTe nanoparticles. Also, the thermal conductivity and heat transfer rates of CdTe nanofluid are greater than that of regular base fluid. Moreover, the thermal conductivity of CdTe-water nanofluid with blade, brick and platelet-shaped particles are 2.40%, 1.40% and 1.91% greater compared with simple water by adding a small amount ($\varphi = 0.005$) of CdTe nanoparticles. The viscosity of platelet-shaped is 11.59% and 17.58% greater than viscosity of blade and brick-shaped nanoparticles, respectively.
Table 8. Thermal conductivity and viscosity of CdTe-water nanofluid for different values of volume fraction $\varphi$.

| $\varphi$ | Thermal Conductivity | Viscosity |
|-----------|----------------------|-----------|
|           | Blade    | Brick    | Platelet | Blade    | Brick    | Platelet |
| 0         | 0.613    | 0.613    | 0.613    | 8.9000×10^{-4} | 8.9000×10^{-4} | 8.9000×10^{-4} |
| 0.005     | 0.6277   | 0.6216   | 0.6247   | 9.5771×10^{-4} | 9.0894×10^{-4} | 10.6873×10^{-4} |
| 0.01      | 0.6425   | 0.6302   | 0.6365   | 10.3091×10^{-4} | 9.4886×10^{-4} | 12.7471×10^{-4} |
| 0.015     | 0.6574   | 0.6389   | 0.6484   | 11.0960×10^{-4} | 10.0916×10^{-4} | 15.0796×10^{-4} |
| 0.02      | 0.6723   | 0.6471   | 0.6604   | 11.9377×10^{-4} | 10.9164×10^{-4} | 17.6847×10^{-4} |
| 0.025     | 0.6874   | 0.6565   | 0.6604   | 12.8344×10^{-4} | 11.9449×10^{-4} | 20.5623×10^{-4} |
| 0.03      | 0.7025   | 0.6654   | 0.6845   | 13.7858×10^{-4} | 13.1832×10^{-4} | 23.5623×10^{-4} |
| 0.035     | 0.7178   | 0.6744   | 0.6967   | 14.7922×10^{-4} | 14.6313×10^{-4} | 27.1355×10^{-4} |
| 0.04      | 0.7331   | 0.6834   | 0.7090   | 15.8534×10^{-4} | 16.2891×10^{-4} | 30.8310×10^{-4} |

6. Conclusions

In this study, a mathematical model of two-dimensional flow of water based nanofluid with suspension of non-spherical CdTe nanoparticles over an inverted cone is analyzed. The thermal conductivity $k_{nf}$ of non-spherical CdTe nanoparticles are obtained by the model in Reference [39]. Magnetic and thermal radiation effects inside a porous medium are also discussed. The effect of active parameters on temperature $T$, velocity $u$, Nusselt number $Nu_x$ and skin friction $\tau_x$ are exhibited via graphs and tables. The analysis showed that

- the blade-shaped particles gained highest Nusselt number compared with platelet and brick type particles. It is reasonable in the sense that blade-shaped particles have higher sphericity than platelet and brick type particles. Therefore, the blade-shaped particles possessed the highest rates for the Nusselt number.
- the temperature distribution increases with an augment in particle volume fraction, radiation parameter and Hartmann number, whereas a reverse relation exists for porosity and Grashof number.
- the highest viscosity is owned by platelet-shaped particles.
- the viscosity of blade-shaped particles is higher than that of brick-shaped particles when ($\varphi < 0.04$).
- the thermal conductivity of blade-shaped particles is 1.93% and 0.94% greater than brick and platelet-shaped nanoparticles when $\varphi = 0.01$.

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Nomenclature

$B_0$ magnetic field strength
$C_p$ heat capacitance
$Gr$ Grashof number
Hartmann number

$K$ dimensionless permeability parameter

$k$ thermal conductivity

$k_0$ permeability of porous medium

$k_b$ absorption coefficient

$L$ reference length

$Nu_x$ Nusselt number

$Pe$ Peclet number

$r$ radius of cone

$Rd$ radiation parameter

$Re$ Reynold number

Greek letters

$\beta$ thermal expansion

$\mu$ dynamic viscosity

$\nabla$ gradient operator

$\nu$ kinematic viscosity

$\phi$ half angle of cone

$\rho$ density

$\sigma$ electrical conductivity

$\sigma_b$ Stefan-Boltzman coefficient

$\tau_s$ skin friction coefficient

$\varphi$ nanoparticle volume fraction

Subscripts

$f$ base fluid

$nf$ nanofluid

$s$ nanoparticles

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