ABSense: Sensing Electromagnetic Waves on Metasurfaces via Ambient Compilation of Full Absorption

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Abstract—Metasurfaces constitute effective media for manipulating and transforming impinging EM waves. Related studies have explored a series of impactful MS capabilities and applications in sectors such as wireless communications, medical imaging and energy harvesting. A key gap in the existing body of work is that the attributes of the EM waves to-be-controlled (e.g., direction, polarity, phase) are known in advance. The present work proposes a practical solution to the EM wave sensing problem using the intelligent and networked MS counterparts—the HyperSurfaces (HSFs), without requiring dedicated field sensors. An nano-network embedded within the HSF iterates over the possible MS configurations, finding the one that fully absorbs the impinging EM wave, hence maximizing the energy distribution within the HSF. Using a distributed consensus approach, the nano-network then matches the found configuration to the most probable EM wave traits, via a static lookup table that can be created during the HSF manufacturing. Realistic simulations demonstrate the potential of the proposed scheme. Moreover, we show that the proposed workflow is the first-of-its-kind embedded EM compiler, i.e., an autonomic HSF that can translate high-level EM behavior objectives to the corresponding, low-level EM actuation commands.

Index Terms—Metasurfaces, HyperSurfaces, EM sensing, nano-networks.

I. INTRODUCTION

Metasurfaces (MS) are highly efficient media for the manipulation of electromagnetic (EM) waves in custom and even unnatural ways. Wave attributes such as the direction of reflection, polarization [1], and the amount of absorption [2], among others, can be controlled with unprecedented accuracy. The outstanding principles MS have been already proven in the microwave, Terahertz (THz), and optical regimes [3], finding their way into a plethora of applications [4]–[6].

The basis for the MS operation is the Huygens principle, which states that any EM wavefront can be traced back to a planar distribution of currents [6]. In that sense, a MS acts as a canvas of currents comprising: i) passive elements acting as sub-wavelength antennas receiving impinging waves, i.e., inductive current sources, and ii) active elements such as PIN diodes acting as current flow manipulators. The active elements can be externally-biased, transforming the currents to follow a planar distribution that corresponds to a required wavefront. Thus, any overall manipulation of the impinging EM waves can be attained, e.g., anomalous steering, polarization and phase alteration, partial attenuation or even full absorption. HSFs are the intelligent, networked variation of MSes, which include an embedded nano-network and a gateway [7]. The nano-network collectively controls the biasing of the active elements in a distributed, autonomic manner, while the gateway connects the nano-network to the HSF-external world via a standard protocol (e.g., WiFi, Bluetooth, etc.).

It is a widespread practice to design MSes (and HSFs) assuming that the impinging wave attributes are known [6]. Nevertheless, this is rarely the case in a broad set of applications where the availability of EM sensing systems becomes a necessity. Currently, the wave attributes can be sensed via HSF-external systems and devices (e.g., d-dot sensors) [8]. However, this approach is not space-granular, since the EM field quantities are sensed i) on average and ii) at one point near or on the HSF gateway. Incorporating multiple such sensors in the HSF adds up to its assembly complexity and overall cost. Moreover, d-EM field sensors are currently size-able, posing novel miniification challenges.

The present study proposes a wave sensing approach that operates without specialized sensory hardware and exploiting the HSF networking capabilities instead. The key idea is to sense the attributes of an impinging wave by fully absorbing it. Full absorption of impinging EM waves is well-studied MS capability [5]. When the surface impedance of the HSF is matched perfectly to the impinging wave, its power dissipates on the passive and active elements, at the same time serving the dual purpose of identifying the condition for perfect absorption. Thus, we use the embedded controller network within the HSF to intelligently iterate over the active element states and obtain the one yielding maximal dissipated power across the active elements. Then, we employ a static lookup table (provided by the HSF manufacturer) that contains the actuator states to achieve full absorption for specific impinging wave cases. The best matching entry is picked as the estimate of the impinging wave attributes. Apart from acting as a wave sensing scheme, the proposed approach is also a form of an ambient EM Compiler [9], since the nano-nodes collectively...
tune the HSF to attain a macroscopic functionality, i.e., the full absorption of an unknown EM wave.

The remainder of the paper is organized as follows. Section II surveys the related studies. Section III details the proposed scheme and evaluation via simulations follows in Section IV. The paper is concluded in Section V.

II. RELATED STUDIES

Metasurfaces are highly effective systems for controlling different aspects of EM waves (wavefront [10], collimation [11], polarization [12], dispersion [13], controllable absorption [14]), especially if they are tunable thus offering the ability to perform multiple functionalities and switch between them at will [3]. In recent years, interest has mainly focused on incorporating voltage-controlled actuation elements inside each unit cell, thus providing a means of dynamically tuning the surface impedance of the metasurface to external EM waves. This is because voltage-controlled elements can be readily controlled by an external controller, such as a Field Programmable Gate Array (FPGA), thus allowing for centrally controlling the metasurface response. The commonly integrated elements are PIN switch [15] or varactor [16] diodes, allowing for a broad range of functionalities ranging from wavefront manipulation, beam splitting, and polarization control. However, switch diodes allow for obtaining two discrete states of the surface impedance; it can be extended to $2^N$ when $N$ unit cells are clustered together to allow for $N$-bit encoding, but $N$ is typically limited to $N = 2$ or $3$ [15] since the resulting supercell must be subwavelength to remain in the metasurface regime. On the other hand, varactor diodes allow for continuous control but only over the reactive part of the surface impedance, thus providing the ability to tune the phase of the impinging wave but not its amplitude.

Recently, it has been shown that having the ability to continuously control both the reactive and resistive part of the complex surface impedance leads to maximum freedom over the available functionalities [17]. This can be achieved by integrating elaborate controller chips in each unit cell that can provide a complex input impedance ($R + jX$). Further enhancing this concept by forming a nano-network of the controllers defines the HyperSurface (HSF) paradigm [7], [18]. In a nutshell, the HyperSurface entails the passive metatom enhanced with an actuation module, a computation and communication module for exchanging data with other cells, and a gateway for communicating with HSF-external entities, such as receiving cell actuation commands and diffusing them for propagation within the inter-cell network. A prominent example of the applications enabled by HSFs are employing mitigation strategies for fault tolerance [19] or deploying programmable wireless environments [5], [18], [20], [21].

Here, we demonstrate a different possibility enabled by the rich HSF capabilities. Specifically, we demonstrate sensing of the characteristics, i.e., direction of incidence (measured by the spherical coordinate angles $\phi$ and $\theta$) and polarization (TE or TM), of the impinging wave. Similar functionalities have been demonstrated in plasmonic metasurfaces incorporating pixelated photodetectors strictly for polarization detection [22] and for orbital angular momentum detection enabled through the decomposition of the transmitted waves impinging on a properly designed metasurface [23]. In our approach we utilize the controllers integrated in the HSF configuration. In particular, we rely on the fact that by fully absorbing the incident wave the power dissipated on the controllers is maximized. More specifically, we monitor the dissipated power on the actuation elements themselves, thus avoiding extra sensing circuitry which would increase the complexity. Moreover, the proposed sensory system is autonomic, not requiring HSF-external computing or communication elements [8]. Additionally, to best of the authors’ knowledge, the present work also constitutes the first application of distributed consensus algorithms [24], [25] as the enabler for the autonomic and collective operation of nano-networks. In this aspect, we note that related works have studied controlled flooding [26], [27], peer-to-peer [28], [29]–[31] and nature-inspired alternatives [32].

It is noted that impressive further capabilities of metasurfaces has been recently demonstrated, promising novel expansions of the HyperSurface concept. Estakhri et al have showcased metastructure design processes able to serve as analog solvers to integral problems [33]. La Spada et al have proposed design workflows for curvilinear metasurfaces exerting arbitrary control over surface EM currents [34], as well as near-zero index wires [35]. Finally, graphene has also been extensively studied for frequency-tunable THz metasurfaces, since its conductivity can be dynamically modulated via electrical, magnetic and optical means [36].
Input 1: An impinging planar wave \([\phi, \theta, \text{polarization}]\); Input 2: The read-only HSF HashMap \(L\): [Parameterized EM Function ID] \(\rightarrow\) [Local Impedance \(z\)]; Output: The sensed wave attributes \([\phi^*, \theta^*, \text{polarization}^*]\); - initialization;

for each mapped value \(z\) of HashMap \(L\) do
- Set all active elements of the HSF to \(z\);
- Measure the power flow \(P_e(z)\) at the active elements of each node \(e\);
end
- Pick \(z_e^* = \arg\max_P\{P_e(z)\}\);
- Obtain the average \(E_e[z_e^*]\) over all nodes;
- Pick \(f_{id}^* = \arg\max_L\{L : z = z_e^*\}\);
- Return parameters \([\phi^*, \theta^*, \text{polarization}^*]\) of \(f_{id}^*\);

Algorithm 1: The workflow of the proposed ABSense scheme.

### III. The Proposed Scheme

We consider a generic HSF architecture as shown in Fig. 1. From its physical aspect, it comprises the common passive and active elements of a metasurface. Additionally, it includes a wireless nano-network embedded within the HSF material, with each nano-node being responsible for the control of one active element. Notably, the HSF material also constitutes the wireless channel medium for the communicating nano-nodes [37, 38]. A distributed communications protocol is employed by the nodes at the application layer, as described later in this Section. The nano-network is supplied with power via an energy pulse generator, which also serves as a rudimentary clock as discussed later. The power pulses can be of any spectrum, such as microwaves, visible light, etc. [39]. A planar wave impinges upon the HSF, with a specific direction and polarization, as illustrated in Fig. 1. The objective is to employ the nano-network capabilities in order to detect these characteristics of the wave, i.e., direction and polarization. The proposed process executed by the HSF as a whole is denoted as ABSense and formulated in pseudo-code as Algorithm 1.

The goal of the ABSense scheme is to sense the impinging wave attributes by: i) detecting the active element configurations that lead to the full absorption of the wave, and ii) deducing the best matching wave attributes by reversing a hash map provided by the HSF manufacturer. This hash map, denoted as \(L\), corresponds any parameterized EM function supported by the HSF to the active element state per node \(e\) to achieve it. For instance, an entry of \(L\) can be written as:

\[
L = \{ f_{id}^{\text{FULL\_ABSORB}}(\phi; 30^\circ, \theta; 42^\circ, TE) \rightarrow z_e : R_1 + jX_1 \}
\]

Treating \(L\) and the impinging wave as inputs, the ABSense workflow is as follows. Signaled and synchronized by the pulse generator, each nano-node iterates over its possible active element states. All nodes move along their iterations in lock-step (detailed below), resulting each time into a uniform surface impedance across the HSF, which will eventually match the impinging wave. For each state \(z\), a node \(e\) obtains a corresponding measurement of power flowing via its active element, \(P_e(z)\). Susceptibility to noise and errors is taken into account. The \(z_e^*\) state that yields the maximum \(P_e\) is communicated to the nano-network as a whole via a distributed consensus approach, and an average value, \(E_e(z_e^*)\), is obtained. Finally, using the \(L\) map in reverse, each nano-node obtains the best matching parameterized function and, hence, the best matching wave attributes as well.

![Fig. 2. The nano-node memory structures and operation in state-chart form (measurements and consensus sub-processes included).](image)

We proceed to detail the operation of each nano-node in Fig. 2. It comprises two phases, the iteration over local active element states (a set of impedance values \(\{z_e : R + jX\}\)) to detect the one \((z_e^*)\) yielding the maximum power flow, and the consensus phase to obtain the \(E_e(z_e^*)\). From another point of view, this phase is also an ambient compilation of an EM functionality, since the nano-nodes self-tune the HSF to perform a full EM wave absorption [9]. The iteration phase is initialized by resetting the \(z_e^*\) variable. Subsequently, the nodes are iterating over all their impedance values, orchestrated by the pulse/power generator (considered to be physically decoupled from the inter-naninode packet exchange process). We assume that the pulse generator emits its energy in the form of pulse sequences representing integer identifiers of the nodes’ impedance values. This pulse emission and orchestration via impedance value identifier broadcast process is continuous. The identifiers are broadcast given enough time for the impedance value to be setup and the power measurement to be obtained (TTM: time to measure), accounting for clock drifts and signal processing variations.

The consensus phase then operates as follows. Each node initializes its consensus value as \(E_e(z_e^*) \rightarrow z_e^* : R + jX\), and broadcasts it to all nano-nodes in its vicinity. A random
delay (RD) is considered to minimize packet collisions. The broadcast date has the form of a packet structured as follows: 

\[
\begin{array}{c|c|c}
\text{SENDER ID (8 bit)} & \text{R (32 bits)} & \text{X (32 bits)}
\end{array}
\]

where the sender ID is an integer identifier of the sending node. During a fixed time interval (TTW: time to wait), each receiving node collects incoming packets while also keeping a log of the average reception power for each one. This information is kept in a hash-map, \( M_{\text{SENDER ID}} \rightarrow \hat{E}[z_{\text{SENDER ID}}]^* \), with sender IDs as keys, as shown in Fig. 2. Subsequently, each node obtains a first estimate of \( E_c[z^*] \) as:

\[
E_c[z^*] \leftarrow E_c[z^*] \cdot w_e + \sum_{\text{SENDER ID}} M_{\text{SENDER ID}} \cdot w_{\text{SENDER ID}}
\tag{2}
\]

where \( w_e \in (0,1) \) is the personal weight that the node gives to its local value, and \( w_{\text{SENDER ID}} \) is the weight assigned to the incoming consensus packets. We consider that the \( w_{\text{SENDER ID}} \) values are proportional to the reception power of the corresponding incoming packets, and are normalized to comply to the condition:

\[
w_e + \sum_{\text{SENDER ID}} w_{\text{SENDER ID}} = 1
\tag{3}
\]

The consensus process converges iteratively, allowing each node to obtain the actual \( E_c[z^*] \) value \[25\]. The consensus process is allowed to run for a maximum number of send/receive packet cycles (max cycles), upon which it yields the estimated \( E_c[z^*] \). The process then concludes by reversing \( L \) to detect the full EM absorption parameterized function that best matches the estimated \( E_c[z^*] \) and, subsequently, it returns the corresponding EM function parameters as the most probable EM wave attributes.

The process as a whole can be immediately restarted, should the HSF tile be operating in a dedicated sensory role, i.e., to sense EM waves and inform other HSF units to adapt accordingly. Alternatively, having obtained the sensory information, the same HSF tile can autonomically apply a different EM manipulation function, such as steering the sensed wave towards a direction. The mode of operation can be adapted to the application scenario at hand.

IV. EVALUATION

We evaluate the proposed scheme using simulations. The evaluation combines full wave simulations conducted in CST \[40\] and nano-network simulations conducted in the AnyLogic platform \[41\]. The full wave simulations produce a dataset of power flow values at each active element in a specific HSF design (shown in Fig. 3) for a variety of impinging wave directions and polarizations, and for any active element state each time. This dataset is then passed on to the nano-network simulator which follows the process described in Section [III].

The HSF used in this work consists of an array of 30 × 30 square metallic patches over a thin metal-backed dielectric (Fig. 3 illustrates a 3 × 3 sub-part for ease of presentation). The unit-cell includes two lumped elements modeled as complex-valued \( RC \) (resistive and capacitive) loads that connect neighbouring metallic patches, along the \( x \)- and \( y \)-directions. The corresponding unit-cell, framed with a white border, is designed for 5 GHz operation and its main parameters are: \( w_{uc} = 10 \) mm, \( w_p = 4.5 \) mm, \( g = 0.5 \) mm, \( t_m = 0.02 \) mm, \( t_d = 0.5 \) mm, \( \varepsilon_r = 2.2 \) and \( \sigma = 5.8 \times 10^7 \) S/m. For normal incidence, perfect absorption of \( x \)- and \( y \)-polarized plane waves is achieved when \( R_x = R_y \approx 1.15 \) Ohm and \( C_x = C_y \approx 0.99 \) pF.

Sample visual hash-maps that are used to link the \( RC_{xy} \) values required for perfect absorption with incidence direction and polarization are depicted in Fig. [4] panel (a) shows the required \( RC_{yz} \) for TE/TM polarized incidence as the elevation angle \( \theta \) varies, for azimuth angle \( \varphi = 0 \), i.e., when \( xz \) is the plane of incidence; panels (b)-(d) show the absorption coefficient for waves of three different incidence directions and polarizations as the \( RC \) values are varied; the white circles mark the regions where the attenuation coefficient is higher than 0.9 (corresponding to reflection coefficient of \(-10 \) dB or lower), and the white crosses mark the \( RC \) values leading to perfect absorption. Note that only the lumped elements oriented parallel to the impinging \( E \)-field affect the resonance of the unit-cell; for instance, when \( \theta, \varphi = (57^o, 0) \) the TE (perpendicular) polarization is only affected by \( RC_y \) elements and TM (parallel) polarization is only affected by \( RC_z \) elements. Finally, our algorithm inherently assumes that the power in the TE and TM polarizations is known (or can be measured) so that the absorption coefficient can be translated to absorbed power. In the context of the consensus workflow, we will assume that the possible \( RC_{yz} \) values are discretized in a \( 10 \times 10 \) grid covering uniformly the axes span of Fig. [4].

The present physical implementation of the HSF can AB-Sense the direction of any linear polarization in the principal planes (\( xz \) and \( yz \), \( \varphi = 0 \) or \( 90^o \), respectively), where the \( x \)- and \( y \)-oriented lumped loads are decoupled and directly correspond to TE or TM polarizations; AB-Sensing the direction of incidence of pure TE or TM (but not both) polarizations, when \( \varphi \in (0, 90^o) \) is also possible, but requires for more careful cross-polarization coupling considerations.
Resistance (Ohm) cycles to cycles value. Finally, as shown cycle value that eliminates the error via the consensus (d) cycles times the RD and TTW). At an

Capacitance (pF)

1.25
1.25
0.8
1.5
0.8
1.5
1
1
(a)
(c)

(b) TE, θ=75°

θ=75°

o

o

0 1 2 3 4
0 1 2 3 4

0 1 2 3 4

Resistance (Ohm) Resistance (Ohm) Capacitance (pF) Capacitance (pF)

(a) RC loads required for perfect absorption of obliquely incident TE and TM polarized waves in the \(xz\) or \(yz\) planes, at 5 GHz and for \(\theta = 0\) to 75° with 5° steps. (b)-(d) Absorption coefficients for a few polarization and incidence direction combinations in the \(xz\) plane, as the \(RC_y\) or \(RC_x\) are varied for TE or TM polarizations, respectively.

![Figure 4](image_url)

as both loads (\(RC_{x/y}\)) affect both polarizations (TE/TM) simultaneously. The general case of ABSensing elliptical polarizations in a broad frequency range and in the entire hemisphere, \(\varphi \in (0, 360^\circ)\), is a topic for future work requiring for more complicated ‘anisotropic’ unit-cell designs (optimized for higher resolution in \(\varphi\theta\) vs. \(RC\) measurement) and, additionally, the ability to measure the current (or voltage) phase on the lumped elements.

Regarding the nano-node workflow parameters, we consider one nano-node per HSF active element. Each nano-node is located exactly below the corresponding element center and at \(t_{z/2}\) depth within the substrate. In order to simulate the inter-nanonode communication channel, we employ the model of \([27, 42, 43]\) using the same physical-layer parameter values (frequency 100 GHz, noise level 0 dBnW, SINR threshold \(\sim 10\) dB, guard interval 0.1 nsec). Assuming a bitrate of 100 Gbps, we consider a consensus data packet duration of approximately 1 nsec (i.e., the 72 bits of the consensus packet plus preamble overheads rounding up at 100 bits total). The transmission power is set to 30 dBnW, yielding approximately 20 nodes within connectivity range.

Regarding the consensus process parameters, we set a TTM equal to 50 times the HSF operating period (5 GHz \(\rightarrow 0.2\) nsec), to accommodate any transient EM phenomena (typically lasting 2 – 3 periods) and obtain dependable average power flow values. The RD is picked at random in the range 0 to 10 packet duration(s) to minimize collisions. The TTW is set to a marginally larger value than the maximum RD, i.e., 12 times the single packet duration. Finally, we assign a 50 % consensus weight to the local optimum of each node and an equal, 50 % aggregate weight to all incoming consensus values.

In Fig. 5 we consider the impinging EM wave of Fig. 4(b).

Additionally, we assume that the power measurements of each node contain a fault expressed as a random percentage over the actual value. We are interested in deducing the minimal max_cycle value that eliminates the error via the consensus value averaging. The consensus process is shown to be very robust to such errors in the measurements. For errors up to \(\sim 82\%\), even one consensus cycle is enough, while \(\sim 12\) cycles can eliminate even the highest measurement errors.

In Fig. 6-7 we focus on the 90 % measurement error case and set max_cycles to 10. Fig. 6 plots the consensus (i.e., sensory) value progression versus time. For ease of exposition, we employ the EM function identifier (cf. rel. (1)) to denote the correct impinging wave attribute (\(f_{id} = 10\)). As shown, the consensus process converges rapidly even for nodes with completely erroneous initial measurement. Thus, the consensus process is economic in terms of required packet transmissions. The corresponding packet statistics are shown in Fig. 7-top: each node received 10 packets (per cycle) from each of its \(\sim 20\) neighbors, subject to some losses due to collisions. The transmission per node are strictly bounded and fully defined by the max_cycles value. Finally, as shown in Fig. 7-bottom, the measurements phase lasts for 10 \(\mu\)sec (i.e., \(10 \times 10\) RC values to iterate over, times the time to obtain a single measurement), while the consensus phase lasts for 1 \(\mu\)sec (i.e., max_cycles times the RD and TTW). At an aggregate EM wave sensing time of 11 \(\mu\)sec, ABSense shows promise for application in real-time-adapting HSFs.
The manipulation and re-shaping of EM waves via intelligent metasurfaces constitutes a key enabler of exotic capabilities in energy harvesting, medical imaging and wireless communications. Nonetheless, accurate EM manipulation requires the accurate sensing of the attributes (direction, polarization, phase) of the wave impinging on a metasurface. Filling in a gap in the related research, the study proposed a novel sensory scheme that exploits the ambient intelligence and communication capabilities of the HyperSurfaces, a novel meta-surface variant. An embedded nano-network automatically tunes the HyperSurface to full absorb the impinging waves, as exhibited by the increase of energy dissipating within it, hence indirectly estimating its attributes. The proposed scheme, validated via extensive EM simulations, can also constitute the basis for distributed EM compiler processes, were the nano-network will auto-tune the HyperSurface status to obtain any high-level objective in its macroscopic EM behavior.

In the future, the authors plan to extend the proposed scheme towards sensing non-planar, complex EM waves, while also advancing the nano-network consensus process into an always-on process that will run in parallel with EM manipulation objectives.

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