On a local systolic inequality for odd-symplectic forms

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Abstract

The aim of this paper is to formulate a local systolic inequality for odd-symplectic forms (also known as Hamiltonian structures) and to establish it in some basic cases.

Let $\Omega$ be an odd-symplectic form on an oriented closed manifold $\Sigma$ of odd dimension. We say that $\Omega$ is Zoll if the trajectories of the flow given by $\Omega$ are the orbits of a free $S^1$-action. After defining the volume of $\Omega$ and the action of its periodic orbits, we prove that the volume and the action satisfy a polynomial equation, provided $\Omega$ is Zoll. This builds the equality case of a conjectural systolic inequality for odd-symplectic forms close to a Zoll one. We prove the conjecture when the $S^1$-action yields a flat $S^1$-bundle or $\Omega$ is quasi-autonomous. In particular the conjecture is established in dimension three.

This new inequality recovers the contact systolic inequality as well as the inequality between the minimal action and the Calabi invariant for Hamiltonian isotopies $C^1$-close to the identity on a closed symplectic manifold. Applications to the study of periodic magnetic geodesics on closed orientable surfaces is given in the companion paper [BK19b].

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1 Introduction

In this article we formulate a local systolic-diastolic inequality for odd-symplectic forms and establish it in some situations. Before discussing the precise set-up, we recall two inequalities which motivate us to study the local systolic-diastolic inequality for odd-symplectic forms and which arise as extreme cases of this inequality. Throughout the paper, we normalise the circle $S^1$ to have length one, namely $S^1 := \mathbb{R}/\mathbb{Z}$.

1.1 Two known inequalities

Contact systolic-diastolic inequality

Let $\Sigma$ be a connected closed manifold of dimension $2n + 1$. Let $\alpha$ be a contact form on $\Sigma$, namely $\alpha$ is a one-form such that $\alpha \wedge (d\alpha)^n$ is nowhere vanishing. Orienting $\Sigma$ so that $\alpha \wedge (d\alpha)^n$ is positive, we define the total volume of $\Sigma$ with respect to $\alpha$ by

$$\text{Volume}(\alpha) := \int_{\Sigma} \alpha \wedge (d\alpha)^n > 0.$$  \hfill (1.1)

We consider the Reeb vector field $R_\alpha$ on $\Sigma$ characterised by the relations $d\alpha(R_\alpha, \cdot) = 0$ and $\alpha(R_\alpha) = 1$ and denote by $\Phi_\alpha$ the generated flow.

A contact form $\alpha_*$ on $\Sigma$ is called Zoll if all orbits of $\Phi_\alpha$ are periodic with the same prime period $T(\alpha_*)$. This is equivalent to saying that $\Phi_\alpha$ induces a free circle action on $\Sigma$ with period $T(\alpha_*)$. In this case, the quotient is a closed symplectic manifold $(M, \omega)$ and the quotient map $p : \Sigma \to M$ is a non-trivial circle bundle. More precisely, we have $p^*\omega = d\alpha_*$ so that $-\frac{1}{T(\alpha_*)} \omega$ represent the Euler class of the bundle. Vice versa, for every closed symplectic manifold $(M, \omega)$ such that $\frac{1}{T} \omega$ represents an integer class for some $T > 0$, there exists a manifold $\Sigma$ and a Zoll contact form with period $T$ yielding $(M, \omega)$ according to the construction above, see [BW58].

By [Bot80], if $\mathfrak{h}$ denotes the free homotopy class of prime periodic orbits of $\Phi_{\alpha_*}$, then $\Phi_\alpha$ carries a periodic orbit in the class $\mathfrak{h}$, if $\alpha$ is sufficiently close to $\alpha_*$. We denote by $T_{\text{min}}(\alpha, \mathfrak{h})$, respectively $T_{\text{max}}(\alpha, \mathfrak{h})$ the minimal, respectively maximal, period of prime periodic orbits of $\Phi_\alpha$ in the class $\mathfrak{h}$. If $\Sigma$ has dimension three, $T_{\text{min}}(\alpha, \mathfrak{h})$ and $T_{\text{max}}(\alpha, \mathfrak{h})$ satisfy the following contact systolic-diastolic inequality, which was originally conjectured by Álvarez-Paiva and Balacheff [APB].
Theorem 1.1. [BK19a, Theorem 1.4] Let \( \Sigma \) be a closed manifold of dimension three, and let \( |H^1_{\text{tor}}(\Sigma; \mathbb{Z})| \) be the order of the torsion part of \( H_1(\Sigma, \mathbb{Z}) \). If \( \alpha_* \) is a Zoll contact form on \( \Sigma \), there exists a \( C^2 \)-neighbourhood \( U \) of \( d\alpha_* \) such that for every contact form \( \alpha \) on \( \Sigma \) with \( d\alpha \in U \), there holds

\[
T_{\min}(\alpha, h)^2 \leq \frac{1}{|H^1_{\text{tor}}(\Sigma; \mathbb{Z})|} \text{Volume}(\alpha) \leq T_{\max}(\alpha, h)^2,
\]

where any of the two equalities holds if and only if \( \alpha \) is Zoll.

Remark 1.2. When \( \Sigma = S^3 \) the result above is due to Abbondandolo, Bramham, Hryniewicz and Salomão, see [ABHS18]. In that paper, forms \( \alpha \) are also constructed, which are \( C^0 \)-close to \( \alpha_* \) and violate the systolic inequality. Furthermore, the same authors showed in [ABHS17] that, if \( (\Sigma, \xi) \) is an arbitrary contact closed three-manifold, then, for every \( C > 0 \), there exists a contact form \( \alpha_C \) with \( \ker \alpha_C = \xi \) such that

\[
T_{\min}(\alpha_C)^2 > C \cdot \text{Volume}(\alpha_C).
\]

Symplectic systolic-diastolic inequality

Let \( (M, \omega) \) be a connected closed symplectic manifold of dimension \( 2n \), and let \( \varphi : M \to M \) be a Hamiltonian diffeomorphism generated by a Hamiltonian \( H : M \times [0, 1] \to \mathbb{R} \). We orient \( M \) so that \( \omega^n \) is positive. One can define the Hamiltonian action \( A_H : \Lambda_{\text{short}}(M) \to \mathbb{R} \) on the space of “short” one-periodic curves \( q : S^1 \to M \) by

\[
A_H(\gamma) := \int_{D^2} \hat{q}^* \omega + \int_0^1 H(q(t), t) dt,
\]

where \( \hat{q} : D^2 \to M \) is a “small” capping disc for \( q \in \Lambda_{\text{short}}(M) \) (see Section 7 for a precise definition). The minimal and maximal Hamiltonian actions of fixed points of \( \varphi \), whose associated curve is short, are given by

\[
\min A_H := \inf_{\gamma \in \text{Crit } A_H} A_H(\gamma), \quad \max A_H := \sup_{\gamma \in \text{Crit } A_H} A_H(\gamma).
\]

The Calabi invariant of \( H \) with respect to \( \omega \) is defined by

\[
\text{CAL}_\omega(H) := \int_{M \times [0, 1]} \omega^n \wedge dt.
\]

One can easily see that the minimal/maximal actions and the Calabi invariant are related in the following way. It is one particular case of Proposition 7.16.

**Proposition 1.3.** Assume that a Hamiltonian diffeomorphism \( \varphi : (M, \omega) \to (M, \omega) \) is generated by a quasi-autonomous Hamiltonian \( H : M \times [0, 1] \to \mathbb{R} \). For instance, \( \varphi \) is the time-one map of a Hamiltonian isotopy \( C^1 \)-close to the identity. Then, there holds

\[
\min A_H \leq \frac{\text{CAL}_\omega(H)}{\int_M \omega^n} \leq \max A_H,
\]

and any of the two equalities holds if and only if \( H(q, t) = h(t) \) for all \((q, t) \in M \times [0, 1]\) and some function \( h : [0, 1] \to \mathbb{R} \). □
Remark 1.4. Proposition 1.3 is false for general Hamiltonian diffeomorphisms. Indeed, adapting [ABHI18, Proposition 2.28], one can construct a time-one map \( \varphi \) of a Hamiltonian isotopy \( C_0 \)-close to the identity not satisfying the inequality above. A proper way to go beyond the quasi-autonomous case might be to take into account not only the fixed points but all the periodic points of \( \varphi \) as explored by Hutchings in [Hut16, Theorem 1.2].

Remark 1.5. In [BK19a] we saw how the use of a global surface of section reduces Theorem 1.1 to a generalised version of Proposition 1.3 for surfaces with boundary, where \( \omega \) is symplectic in the interior and vanishes of order one at the boundary.

To give a unified model for the above two situations, we notice that the Hamiltonian diffeomorphism \( \varphi : M \to M \) can be seen as the return map of a flow \( \Phi^X \) on the manifold \( M \times S^1 \). If \( H : M \times S^1 \to \mathbb{R} \) is a Hamiltonian one-periodic in time generating \( \varphi \) as time one-map, then \( X(q,t) := \partial_t + X_{H_t}(q) \), for all \( (q,t) \in M \times S^1 \). Here, \( X_{H_t} \) is the Hamiltonian vector field tangent to \( M \times \{ t \} \) associated with the function \( H_t := H(\cdot, t) \). Thus, we view Proposition 1.3 as giving a systolic-diastolic inequality for those flows on the trivial \( S^1 \)-bundle \( p : M \times S^1 \to M \), which are obtained by lifting a Hamiltonian isotopy of \( M \). On the other hand, Theorem 1.1 yields a systolic-diastic inequality for Reeb flows on the non-trivial \( S^1 \)-bundles obtained from the Boothby-Wang construction. We will show that these two types of flows are particular incarnations of the flow induced by an odd-symplectic structure (also known as Hamiltonian structure, see [CM05]) on an oriented odd-dimensional closed manifold which is the total space of an arbitrary oriented circle bundle over a symplectic manifold. Moreover, we will formulate a conjectural systolic-diastolic inequality in this setting which recovers the contact (Theorem 1.1) and the symplectic one (Proposition 1.3) as two extreme cases.

1.2 Odd-symplectic systolic geometry

Let \((\Sigma,\partial\Sigma)\) be a connected oriented closed manifold of dimension \(2n+1\), and let \( C \in H^2_{dR}(\Sigma) \) be a class in its de Rham cohomology. We write \( \Omega^1(\Sigma) \) for the set of one-forms on \( \Sigma \) and \( \Xi^2_C(\Sigma) \) for the set of closed two-forms on \( \Sigma \) representing the class \( C \). We pick an auxiliary element \( \Omega_0 \in \Xi^2_C(\Sigma) \) and get a surjective map

\[
\Omega^1(\Sigma) \to \Xi^2_C(\Sigma), \quad \alpha \mapsto \Omega_\alpha := \Omega_0 + d\alpha, \quad (1.4)
\]

whose kernel is the space of closed one-forms on \( \Sigma \). Using this map, we can associate to \( \alpha \in \Omega^1(\Sigma) \) a real number \( \text{Vol}(\alpha) \), which generalises both the contact volume (1.1) and the Calabi invariant (1.3). Namely, we define a functional \( \text{Vol} : \Omega^1(\Sigma) \to \mathbb{R} \) by

\[
\text{Vol}(0) = 0, \quad d_\alpha \text{Vol} \cdot \beta = \int_\Sigma \beta \wedge \Omega^n_\alpha, \quad \forall \beta \in \Omega^1(\Sigma).
\]

For instance, when \( n = 1 \), we have

\[
\text{Vol}(0) = 0, \quad \text{Vol}(\alpha) = \frac{1}{2} \int_\Sigma \alpha \wedge d\alpha + \int_\Sigma \alpha \wedge \Omega_0.
\]

We also remark that, when \( C = 0 \) and \( \Omega_0 = 0 \), the function \( \text{Vol} \) recovers the Chern-Simons action for principal \( S^1 \)-bundles over \( \Sigma \) (see [CS74] and Remark 2.1). For \( n = 1 \), \( \Omega_0 = 0 \) and
\[ d\alpha = \iota_F \mu \] for some vector field \( F \) on \( \Sigma \), whose flow preserves a global volume form \( \mu \) on \( \Sigma \), \( \text{Vol}(\alpha) \) recovers the helicity of \( F \), see \([\text{Pat09}]\).

If \( C^n = 0 \), then we can use the map \([1,4]\) to push forward \( \text{Vol} \) to a **volume functional**

\[ \mathcal{Vol} : \Xi^2_\mathcal{C}(\Sigma) \rightarrow \mathbb{R}. \]

If \( C^n \neq 0 \), this procedure does not work, since there exists \( \tau \in H^1_{\text{dR}}(\Sigma) \) such that \( \tau \cup C^n \neq 0 \). In this case, we say that \( \alpha \in \Omega^1(\Sigma) \) is **normalised**, if \( \text{Vol}(\alpha) = 0 \). For every \( \Omega \in \Xi^2_\mathcal{C}(\Sigma) \), there exists \( \alpha \in \Omega^1(\Sigma) \) normalised such that \( \Omega = \Omega_\alpha \). Therefore, in this situation we can just work with normalised forms and declare \( \mathcal{Vol} : \Xi^2_\mathcal{C}(\Sigma) \rightarrow \mathbb{R} \) to be identically zero. In both cases, the volume functional is invariant under diffeomorphisms \( \Psi : \Sigma \rightarrow \Sigma \) isotopic to the identity (see Proposition 2.8):

\[ \mathcal{Vol}(\Psi^*\Omega) = \mathcal{Vol}(\Omega). \]

Having identified what the volume should be, we want to introduce the set \( \mathcal{X}(\Omega) \) of closed characteristics of \( \Omega \) generalising periodic Reeb and Hamiltonian orbits. To this purpose, we consider the possibly singular distribution \( \ker \Omega \rightarrow \Sigma \) called the **characteristic distribution** and define

\[ \mathcal{X}(\Omega) := \left\{ \gamma : S^1 \hookrightarrow \Sigma \mid \dot{\gamma} \in \ker \Omega \right\} / \sim, \]

where \( \gamma_1 \sim \gamma_2 \) if and only if \( \gamma_1 \) and \( \gamma_2 \) coincide up to an orientation-preserving reparametrisation of \( S^1 \). The distribution \( \ker \Omega \) is co-oriented by \( \Omega^n \) and we orient it using the given orientation \( o_\Sigma \) on \( \Sigma \).

We single out the forms in \( \Xi^2_\mathcal{C}(\Sigma) \), whose characteristic distribution is one-dimensional. They are of special importance, as their closed characteristics are the periodic orbits of a flow.

**Definition 1.6.** A two-form \( \Omega \) on \( \Sigma \) is said to be **odd-symplectic**, if it is closed and maximally non-degenerate, namely if its characteristic distribution is an (oriented) real line bundle over \( \Sigma \). We write \( \mathcal{S}_C(\Sigma) \) for the subset of all odd-symplectic forms in \( \Xi^2_\mathcal{C}(\Sigma) \).

In \( \mathcal{S}_C(\Sigma) \) we hope to find elements whose characteristic distribution is as simple as possible. To this purpose we introduce the set of oriented \( S^1 \)-bundles with total space \( \Sigma \):

\[ \mathfrak{P}(\Sigma) := \{ p : \Sigma \rightarrow M \mid p \text{ is an oriented } S^1\text{-bundle} \}. \]

**Definition 1.7.** An odd-symplectic form \( \Omega \) is said to be **Zoll** if the oriented leaves of its characteristic distribution are the fibres of some \( p_\Omega : \Sigma \rightarrow M_\Omega \) in \( \mathfrak{P}(\Sigma) \). We write \( \mathcal{Z}_C(\Sigma) \) for the set of Zoll forms in \( \mathcal{S}_C(\Sigma) \).

**Example 1.8.** Let \( (N_1, \sigma_1) \) and \( (N_2, \sigma_2) \) be two connected closed symplectic manifolds and suppose that \( [\sigma_1] \in H^2_{\text{dR}}(N_1) \) is an integral cohomology class. We build the product symplectic manifold \( (N_1 \times N_2, \sigma_1 \oplus \sigma_2) \) and consider the oriented \( S^1 \)-bundle \( p : \Sigma \rightarrow N_1 \times N_2 \) whose Euler class is \(-[\sigma_1 \oplus 0]\). Then, \( \Omega := p^*(\sigma_1 \oplus \sigma_2) \) is a Zoll form with cohomology class \( C := p^*[0 \oplus \sigma_2] \).

Odd-symplectic Zoll forms \( \Omega \) with vanishing cohomology class are exactly the differentials of Zoll contact forms \( \alpha \). We can use this observation to extend the classification of Zoll contact forms on three-manifolds contained in \( [\text{BK19a}, \text{Proposition 1.2}] \) to Zoll odd-symplectic forms.

**Proposition 1.9.** Let \( \Sigma \) be a connected oriented closed three-manifold. There is a Zoll odd-symplectic form on \( \Sigma \) if and only if \( \Sigma \) is the total space of an oriented \( S^1 \)-bundle over a...
connected oriented closed surface $M$. When the bundle is non-trivial, then every Zoll odd-symplectic form is the differential of a Zoll contact form. When the bundle is trivial, then the group $H_1(\Sigma; \mathbb{Z})$ is free with rank equal to $2 \text{genus}(M) + 1$, and if $\Omega$ and $\Omega'$ are Zoll odd-symplectic forms on $\Sigma$, there exists a real number $T > 0$ and a diffeomorphism $\Phi : \Sigma \to \Sigma$ such that $\Phi^* \Omega = T \Omega$. Moreover,

- if $\Sigma = S^2 \times S^1$, the set $\mathcal{Z}(\Sigma)$ has exactly two connected components, $\forall C \in H^2_{\text{dr}}(\Sigma) \setminus 0$;
- if $\Sigma = \mathbb{T}^2 \times S^1$, the set $\mathcal{Z}(\Sigma)$ is non-empty and connected, $\forall C \in H^2_{\text{dr}}(\Sigma) \setminus 0$;
- if $\Sigma = M \times S^1$ with $\chi(M) < 0$, there is a one-dimensional subspace $L \subset H^2_{\text{dr}}(\Sigma)$ such that $\mathcal{Z}(\Sigma)$ is non-empty if and only if $C \in L \setminus 0$; in this case the set $\mathcal{Z}(\Sigma)$ is connected.

**Remark 1.10.** The classification of Zoll odd-symplectic forms up to diffeomorphism on a three-manifold is equivalent to the classification of bundles in the set $\{p_\Omega \mid \Omega \in \mathcal{Z}(\Sigma)\}$ up to isomorphism. Analogously, the connected components of $\mathcal{Z}(\Sigma)$ on a three-manifold are in bijection with the connected components of $\{p_\Omega \mid \Omega \in \mathcal{Z}(\Sigma)\}$. This is due to the fact that the map $\mathcal{Z}(\Sigma) \to \mathcal{F}(\Sigma)$, given by $\Omega \mapsto p_\Omega$, has contractible fibres.

Let us assume that $\mathcal{Z}(\Sigma)$ is not empty and take $\Omega_* \in \mathcal{Z}(\Sigma)$ with associated $S^1$-bundle $p_{\Omega_*} : \Sigma \to M_*$. This implies that there exists a positive symplectic form $\omega_*$ on $M_*$ such that $\Omega_* = p_{\Omega_*}^* \omega_*$. We set $c_* := [\omega_*] \in H^2_{\text{dr}}(M_*)$. We write $\mathfrak{h} \in [S^1, \Sigma]$ for the free-homotopy class of the oriented $p_{\Omega_*}$-fibres and $e_* \in H^2_{\text{dr}}(M_*)$ for minus the real Euler class of $p_{\Omega_*}$. Let $\mathcal{F}(\Sigma)$ be the connected component of $p_{\Omega_*}$ inside $\mathcal{F}(\Sigma)$. As we did for the volume, the action will be computed with respect to some reference object, which we now define.

**Definition 1.11.** A **weakly Zoll pair** is a couple $(p, c)$, where $p : \Sigma \to M$ is an element in $\mathcal{F}(\Sigma)$ and $c \in H^2_{\text{dr}}(M)$ is a cohomology class. We write $\mathcal{F}_c(\Sigma)$ for the set of weakly Zoll pairs and $\mathcal{F}_c(\Sigma)$ for the subset of those pairs $(p, c)$ such that $C = p^* c \in H^2_{\text{dr}}(\Sigma)$. We denote by $\mathcal{F}_c^0(\Sigma)$ the set of those $(p, c) \in \mathcal{F}_c(\Sigma)$ such that $p \in \mathcal{F}_c^0(\Sigma)$.

**Remark 1.12.** If $\Omega \in \mathcal{F}_c(\Sigma)$ is Zoll, then $(p_\Omega, [\omega_\Omega]) \in \mathcal{F}_c(\Sigma)$ is a weakly Zoll pair, where $\omega_\Omega$ is the closed two-form on $M_\Omega$ such that $\Omega = p_\Omega^* \omega_\Omega$. Conversely, if $(p, c) \in \mathcal{F}_c(\Sigma)$ is a weakly Zoll pair, we can consider any closed two-form $\omega$ on $M$ such that $c = [\omega]$ and build the two-form $p^* \omega \in \mathcal{F}_c(\Sigma)$, which is Zoll exactly when $\omega$ is symplectic.

Let us now fix a reference weakly Zoll pair

$$(p_0, c_0) \in \mathcal{F}_c^0(\Sigma), \quad p_0 : \Sigma \to M_0.$$  

We write $e_0 \in H^2_{\text{dr}}(M_0)$ for minus the real Euler class of $p_0$ and we have the equivalence

$$C_0 \neq 0 \iff e_0 = 0.$$  

Let $\mathcal{A}_0(\Sigma)$ be the space of one-periodic curves in the class $\mathfrak{h}$, and let $\tilde{\mathcal{A}}_0(\Sigma)$ be the space of homotopies of paths $\{\gamma_t\}_{t \in [0,1]}$ inside $\mathcal{A}_0(\Sigma)$ such that $\gamma_0$ is some oriented $p_0$-fibre. The admissible homotopies are allowed to move $\gamma_0$ inside the set of oriented $p_0$-fibres but have to fix the periodic curve $\gamma_1$, so that the natural projection $\tilde{\mathcal{A}}_0(\Sigma) \to \mathcal{A}_0(\Sigma)$, $[\gamma_t] \mapsto \gamma_1$ is a covering map.
We pick a closed two-form \( \omega_0 \in \Xi^2(M_0) \) and choose \( \Omega_0 := p_0^* \omega_0 \in \Xi^2_c(\Sigma) \) as our auxiliary element for the computation of the volume. One sees that \( \Psi \otimes : \Xi^2_c(\Sigma) \to \mathbb{R} \) depends only on \((p_0, c_0)\) and not on \( \omega_0 \). We associate to \( \Omega = \Omega_0 + d\alpha \) the action functional

\[
\tilde{A}_\Omega : \Lambda^0(\Sigma) \to \mathbb{R}, \quad \tilde{A}_\Omega([\gamma_r]) := \int_{[0,1] \times S^1} \Gamma^* \Omega + \int_{S^1} \gamma_0^* \alpha,
\]

where \( \alpha \) is chosen to be normalised if \( C^n \neq 0 \), and \( \Gamma : [0,1] \times S^1 \) is the cylinder traced by the path \( \{ \gamma_r \} \). Like the volume, the functional \( \tilde{A}_\Omega \) depends only on \((p_0, c_0)\) and not on \( \omega_0 \). If \( [\gamma_r] \) is a critical point of \( \tilde{A}_\Omega \), then \( [\gamma] \in {\mathcal{A}}(\Omega) \), provided \( \gamma \) is embedded. Furthermore, the action is invariant under isotopies \( \{ \Psi_r : \Sigma \to \Sigma \} \) starting at the identity (Proposition 6.10):

\[
\tilde{A}_{\Psi_1^* \Omega}([\Psi_r^{-1} \circ \gamma_r]) = \tilde{A}_\Omega([\gamma_r]), \quad \forall [\gamma_r] \in \Lambda^0(\Sigma).
\]

In general, \( \tilde{A}_\Omega \) does not descend to a functional on \( \Lambda^0(\Sigma) \). More precisely, by Lemma 6.3 this happens if and only if \( c_0|_{\pi_2(M_0)} = ac_0|_{\pi_2(M_0)} \) for some \( a \in \mathbb{R} \). However, as we see now, we can define an action on the set of Zoll forms \( \Omega \) with \((p_0, [\omega_0]) \in \Xi^0_c(\Sigma) \). In this case, the volume of \( \Omega \) can be expressed as a polynomial function of the action.

**Definition 1.13.** The Zoll polynomial \( P : \mathbb{R} \to \mathbb{R} \) associated with \((p_0, c_0)\) is given by

\[ P(0) = 0, \quad \frac{dP}{dA}(A) = \langle (Ac_0 + c_0)^n, [M_0] \rangle. \]

For instance, when \( n = 1 \), the polynomial reads

\[ P(A) = \frac{1}{2} \langle c_0, [M_0] \rangle A^2 + \langle c_0, [M_0] \rangle A. \]

**Theorem 1.14.** There is a well-defined volume function

\[ \Psi \otimes : \Xi^2_c(\Sigma) \to \mathbb{R}, \quad \Psi \otimes(p, c) := \Psi \otimes(p^* \omega) \]

and well-defined action functional

\[ A : \Xi^0_c(\Sigma) \to \mathbb{R}, \quad A(p, c) := \tilde{A}_{p^* \omega}([\delta_r]). \]

Here, \( \omega \) is any closed two-form on \( M \) in the class \( c \), \( \{ \delta_r \} \) is any path of periodic curves such that \( \delta_r \) is an oriented \( p_r \)-fibre, where \( \{ p_r \} \) is any path of oriented \( S^1 \)-bundles from \( p_0 \) to \( p_1 = p \). Moreover, there holds

\[ P(A(p, c)) = \Psi \otimes(p, c), \quad \forall (p, c) \in \Xi^0_c(\Sigma). \]

If \( A_* := A(p_{\Omega_*}, c_*) \), then \( \frac{dP}{dA}(A_*) > 0 \). In particular, the polynomial \( P \) is non-zero.

From this result, we can generalise the equality cases in Theorem 1.1 and Proposition 1.3.

**Corollary 1.15.** Let \( \Omega \in \Xi^2_c(\Sigma) \) be a Zoll odd-symplectic form such that \( p_0 \in \Xi^0_c(\Sigma) \). If we set \( A(\Omega) := A(p_\Omega, [\omega_\Omega]) \), then

\[ P(A(\Omega)) = \Psi \otimes(\Omega). \]
In what follows, we describe a conjectural systolic-diastolic inequality for odd-symplectic forms close to \( \Omega_* \) and with class \( C \in H^2_{dR}(\Sigma) \). To this end, we fix a finite open covering \( \{ B_i \} \) of \( M_* \) by balls so that all their pairwise intersections are also contractible. Let \( \Lambda(\Omega_\alpha) \) be the space of periodic curves \( \gamma \in \Lambda_b(\Sigma) \) with the property that \( \rho_\Omega(\gamma) \) is contained in some \( B_i \) and there is a path \( \{ \gamma^\text{short}_r \} \in \Lambda_b(\Sigma) \) entirely contained in \( p^{-1}_O(B_i) \) with \( \gamma^\text{short}_1 = \gamma \). If \( \Omega \in \mathcal{E}_C(S) \), we set
\[
\mathcal{A}_\Omega : \Lambda(\Omega_\alpha) \to \mathbb{R}, \quad \mathcal{A}_\Omega(\gamma) := \mathcal{A}_O([\delta_r \# \gamma^\text{short}_r]), \quad \gamma^\text{short}_1 = \gamma.
\]
Here, the symbol \# denotes the concatenation of paths, \( \{ \delta_r \} \) is any path of periodic curves such that \( \delta_1 = \gamma^\text{short}_0 \), and for every \( r \in [0,1] \), \( \delta_r \) is an oriented \( p_r \)-fibre, where \( \{ p_r \} \) is a path of oriented \( S^1 \)-bundles connecting \( p_0 \) with \( p_1 = p_{\Omega_*} \). We define
\[
\mathcal{A}_\min(\Omega) := \inf_{\gamma \in \mathcal{X}(\Omega) \cap \Lambda(\rho_\Omega)} \mathcal{A}_\Omega(\gamma), \quad \mathcal{A}_\max(\Omega) := \sup_{\gamma \in \mathcal{X}(\Omega) \cap \Lambda(\rho_\Omega)} \mathcal{A}_\Omega(\gamma).
\]
By [Gins87 Section III] or [APB14 Section 3.2], if \( \Omega \in \mathcal{S}_C(S) \) is \( C^1 \)-close to \( \Omega_* \), the set \( \mathcal{X}(\Omega) \cap \Lambda(\rho_\Omega) \) is compact and non-empty. Furthermore, the numbers \( \mathcal{A}_\min(\Omega) \) and \( \mathcal{A}_\max(\Omega) \) are finite and vary \( C^1 \)-continuously with \( \Omega \).

**Conjecture 1** (Local systolic-diastolic inequality for odd-symplectic forms). Let \( \Omega_* \) be a Zoll odd-symplectic form with cohomology class \( C \in H^2_{dR}(\Sigma) \). There is a \( C^k\)-neighbourhood \( \mathcal{U} \) of \( \Omega_* \) in \( \mathcal{S}_C(S) \) with \( k \geq 2 \) such that
\[
P(\mathcal{A}_\min(\Omega)) \leq \mathfrak{Vol}(\Omega) \leq P(\mathcal{A}_\max(\Omega)), \quad \forall \Omega \in \mathcal{U}.
\]
The equality holds in any of the two inequalities above, if and only if \( \Omega \) is Zoll.

**Remark 1.16.** If the real Euler class of the bundle associated with \( \Omega_* \) vanishes, then the inequality in the conjecture is equivalent to
\[
\mathcal{A}_\min(\Omega) \leq 0 \leq \mathcal{A}_\max(\Omega), \quad \forall \Omega \in \mathcal{U},
\]
with any of the equalities holding if and only if \( \Omega \) is Zoll.

**Remark 1.17.** The volume, the action and the Zoll polynomial depend on the choice of reference pair \( (p_0, c_0) \in \mathfrak{E}^0_C(S) \). However, thanks to Theorem [L14] we will observe in Remark [7.7] that the functional
\[
(P \circ \mathcal{A}_\Omega - \mathfrak{Vol}(\Omega)) : \Lambda(\rho_{\Omega_*}) \to \mathbb{R}
\]
is independent of such a choice. Hence, Conjecture [H] remains the same when we use another reference pair in \( \mathfrak{E}^0_C(S) \).

**Remark 1.18.** As mentioned at the beginning of this subsection, Conjecture [H] recovers the contact and symplectic systolic-diastolic inequality. In the contact case, we have \( \Omega_* = d\alpha_* \) for a contact form \( \alpha_* \) so that \( C = 0 \). Up to multiplication by a positive constant, we can assume that \( T(\alpha_*) = 1 \). We take \( \varrho_\Sigma = \varrho_{\alpha_*} \) and \( c_0 = 0 \). There holds
\[
P(A) = \frac{1}{n + 1} \langle e^n_0, [M_0] \rangle A^{n+1},
\]
where \( \langle e^n_0, [M_0] \rangle = \langle [\omega^n_0], [M_0] \rangle > 0 \). When \( n = 1 \), we have \( \langle e_0, [M_0] \rangle = |H^1_{tor}(\Sigma; \mathbb{Z})| \). If \( \Omega = d\alpha \) is \( C^0 \)-close to \( \Omega_* \), we can take \( \alpha \) to be a contact form. We have
\[
\mathfrak{Vol}(d\alpha) = \frac{1}{n} \text{Volume}(\alpha), \quad A_{d\alpha}(\gamma) = \int_{S^1} \gamma^* \alpha = T(\gamma), \quad \forall \gamma \in \Lambda_0(\Sigma).
\]
Thus, there holds $T_{\min}(\alpha, h) \leq A_{\min}(d\alpha) \leq A_{\max}(d\alpha) \leq T_{\max}(\alpha, h)$, and Conjecture 1 implies
\[
T_{\min}(\alpha, h)^n \leq \frac{1}{\langle e_0^n, [M_0]\rangle} \text{Volume}(\alpha) \leq T_{\max}(\alpha, h)^n
\]
with equality cases exactly when the contact form $\alpha$ is Zoll.

In the symplectic case, $p_\Omega$ is trivial so that $\Sigma = M_s \times S^1$ and we call $t \in S^1$ the global angular coordinate. We take $\Omega = \Omega_\ast + d(Hdt)$ for some function $H : M_s \times S^1 \to \mathbb{R}$. The characteristic distribution of $\Omega$ is generated by the vector field $\partial_t + X_H$, where $X_H$ is the $\omega_\ast$-Hamiltonian vector field of $H_t := H(\cdot, t)$ tangent to $M_s \times \{t\}$. Then, $A_\Omega$ is the Hamiltonian action defined in (1.2), $\langle e_0^n, [M_0]\rangle = \langle [\omega^n_0], [M_0]\rangle > 0$, and
\[
P(A) = \langle e_0^n, [M_0]\rangle A, \quad \text{Vol}(Hdt) = \int_{M_s \times S^1} (Hdt) \wedge \omega^n_\ast = \text{CAL}_{\omega_\ast}(H).
\]
Conjecture 1 generalises Proposition 1.3 since $A_{\min}(\Omega) = \min A_H$ and $A_{\max}(\Omega) = \max A_H$, if we take the Hamiltonian $H$ to be normalised, namely $\text{CAL}_{\omega_\ast}(H) = 0$.

Inspired by the Hamiltonian case, we will look first at a special class of $\Omega \in \Xi^2_C(\Sigma)$ when studying Conjecture 1. Here, we assume without loss of generality that $p_0 = p_{\Omega_\ast}$. We fix an $S^1$-connection form $\eta$ for the bundle $p_0$. This means that $\eta$ restricts to the angular form on each fibre and $d\eta = p_0^* \kappa$ for some $\kappa \in \Xi^2_C(M_0)$. We consequently choose the form $\omega_\ast$ so that $\omega_\ast = \omega_0 + A_\ast \kappa$, where $A_\ast = A(\Omega_\ast)$. We say that a form $\Omega \in \Xi^2_C(\Sigma)$ is an $H$-form, if there exists a function $H : \Sigma \to \mathbb{R}$ such that
\[
\Omega = \Omega_0 + d(H\eta).
\]
We call $H$ a defining Hamiltonian for $\Omega$. If $H$ is $C^0$-close to $A_\ast$, then $\Omega$ is odd-symplectic, and if $H$ is $C^k$-close to $A_\ast$, then $\Omega$ is $C^{k-1}$-close to $\Omega_\ast$. An $H$-form is called quasi-autonomous if there exist $q_{\min}, q_{\max} \in M_0$ and defining Hamiltonians $H_{\min}$ and $H_{\max}$ such that
\[
\min_\Sigma H_{\min} = H_{\min}(z), \quad \forall \, z \in p_0^{-1}(q_{\min}), \quad \max_\Sigma H_{\max} = H_{\max}(z), \quad \forall \, z \in p_0^{-1}(q_{\max}).
\]
We say that $H$ is quasi-autonomous, if the corresponding $\Omega$ is quasi-autonomous. We establish Conjecture 1 for quasi-autonomous forms close to a Zoll odd-symplectic form.

**Proposition 1.19.** There exists a $C^2$-neighbourhood $\mathcal{H}$ of the constant $A_\ast$ in the space of quasi-autonomous functions over $\Sigma$ such that if $H \in \mathcal{H}$ and $\Omega = \Omega_\ast + d(H\eta)$ then
\[
P(A_{\min}(\Omega)) \leq \mathfrak{Sol}(\Omega) \leq P(A_{\max}(\Omega)),
\]
and any of the two equalities holds if and only if $\Omega$ is Zoll. In the Zoll case, $H$ is invariant under the holonomy of $\eta$. In particular, it is constant if $e_0 \neq 0$.

Let $p_0 : M_0 \times S^1 \to M_0$ be the trivial bundle and take $\eta = dt$, where $t \in S^1$ is the angular coordinate. In this case, we show in Proposition 1.3 that if $\Omega \in \Xi^2_C(\Sigma)$ is $C^2$-close to $\Omega_\ast$, then $\Omega$ is an $H$-form with a $C^2$-small defining Hamiltonian $H$, after applying a diffeomorphism of $\Sigma$ isotopic to the identity. Therefore, in this setting the conjecture follows from Proposition 1.3 (or Proposition 1.19) and the invariance of the volume and the action under diffeomorphisms. From this result and the topological Lemma 4.10 Conjecture 1 for bundles with $e_0 = 0$ can readily be proven.
Theorem 1.20. The local systolic-diastolic inequality for odd-symplectic forms, whose associated bundle has vanishing real Euler class, holds true in the $C^2$-topology.

If $\Omega$ is any Zoll odd-symplectic form on a closed three-manifold $\Sigma$, then either the Euler class of its associated bundle vanishes or $\Omega = d\alpha$ for some Zoll contact form $\alpha$. Hence, Theorem 1.20 together with [BK19a, Theorem 1.4] establish Conjecture 1 in dimension three.

Corollary 1.21. The local systolic-diastolic inequality for odd-symplectic forms holds true in the $C^2$-topology on closed three-manifolds.

Remark 1.22. More generally, one could try to formulate a systolic-diastolic inequality in a neighbourhood of an odd-symplectic form $\Omega$, whose closed characteristics are tangent to a (not necessarily free) $S^1$-action on $\Sigma$, cf. [Tho76] and [BR94].

The example of magnetic geodesics

Our main motivation to study the systolic-diastolic inequality for odd-symplectic forms comes from twisted cotangent bundles, where the magnetic form is symplectic. Let $(N, \sigma)$ be a connected closed symplectic manifold of dimension $2m$. Let $p_{T^*N} : T^*N \to N$ be the cotangent bundle map and consider the twisted symplectic form $d\lambda_{\text{can}} + p^*_{T^*N}\sigma \in \Omega^2(T^*N)$,

where $\lambda_{\text{can}}$ is the canonical one-form on $T^*N$. We fix a $\sigma$-compatible almost complex structure $J$ on $N$ with associated metric $g_J$. The structure $J$ turns $p_{T^*N}$ into a complex vector bundle and we denote by $p_g : S^*_gN \to N$ and $i_g : S^*_gN \hookrightarrow T^*N$ for the natural projection and inclusion. The one-parameter group $t \mapsto e^{tJ}$ acts fibrewise on $T^*N$ and yields a free $S^1$-action on $S^*_gN$, since $J$ is $g$-orthogonal. The quotient is naturally identified with $\mathbb{P}_C(T^*N)$, so that we have an oriented $S^1$-bundle $p$ making the following diagram commute

\[
\begin{array}{ccc}
T^*N & \overset{i_g}{\longrightarrow} & S^*_gN \\
p_{T^*N} \downarrow & & \downarrow p_g \\
N & \underset{\mathbb{P}_C(T^*N)}{\longrightarrow} & \mathbb{P}_C(T^*N).
\end{array}
\]

Therefore, $\Omega_{g,\sigma} := i_g^*(d\lambda_{\text{can}} + p^*_{T^*N}\sigma)$ is an odd-symplectic two-form on $S^*_gN$ with cohomology class $C := p_g^*[\sigma]$. The class $C$ vanishes if and only if $N$ is a surface different from the two-torus. In general, it is well-known that there exists a Zoll form $p^*\omega_\sigma$ in the class $C$. More specifically, $\omega_\sigma$ is the symplectic form on $\mathbb{P}_C(T^*N)$ defined as

\[
\omega_\sigma := a\omega_\text{FS} + p^*J\sigma,
\]

for some $a > 0$ small enough (see [Voi07, Proposition 3.18] when $(N, \sigma, J)$ is Kähler). Here, $p^*\omega_{\text{FS}} = d\eta$, where $\eta$ is a connection form for $p$, which, for all $x \in N$, restricts on the fibre $p^{-1}_g(x) \cong S^{2m-1}$ to the standard contact form on the sphere. The $\mathbb{P}$-fibres are almost tangent to the characteristic distribution of $\Omega_{g,\sigma}$, if $\sigma$ is very big. However, the form $\Omega_{g,\sigma}$ and $p^*\omega_\sigma$ are remarkably not close to each other, if $m > 1$. 
The relevance of this example stems from the fact that the characteristics of $\Omega_{g,\sigma}$ are the tangent lifts of the magnetic geodesics of the pair $(g,\sigma)$. A curve $c : \mathbb{R} \to N$ is called a magnetic geodesic if $g(\dot{c},\dot{c}) \equiv 1$ and there holds

$$\nabla^g_{\dot{c}} \dot{c} = -fJ\dot{c},$$

where $\nabla^g$ is the Levi-Civita connection of $g$ and $f : N \to (0,\infty)$ is the conformal factor given by $f \cdot g = g_J$. We refer to the work of Kerman in [Ker99] for results on the existence of closed magnetic geodesics in this setting.

In the companion paper [BK19b], we use Corollary 1.21 and Proposition 1.9 to establish a systolic-diastolic inequality for magnetic geodesics when $N$ is a surface, namely $m = 1$.

1.3 Structure of the paper

- In Section 2, we introduce the volume of a closed two-form on an odd-dimensional oriented closed manifold and explore its invariance properties.

- We define odd-symplectic forms and H-forms in Section 3. Under certain conditions, we prove a stability result for H-forms in the set of odd-symplectic forms.

- We devote Section 4 to review some basic facts about oriented $S^1$-bundles, which will be crucially used in the following discussion.

- In Section 5, we define weakly Zoll pairs and Zoll odd-symplectic forms. We also prove Proposition 1.9 which classifies Zoll odd-symplectic forms on oriented three-manifolds.

- In Section 6, we introduce the action of a closed two-form on an odd-dimensional manifold which is the total space of some oriented $S^1$-bundle. We prove Theorem 6.14, which is just a reformulation of Theorem 1.14 showing that the action and the volume of a weakly Zoll pair are related through the Zoll polynomial.

- We formulate the local systolic-diastolic inequality for odd-symplectic forms in Section 7. We prove it for quasi-autonomous H-forms (Proposition 7.16 corresponding to Proposition 1.19) or when the real Euler class of the bundle vanishes (end of Subsection 7.3 corresponding to Theorem 1.20 and Proposition 1.3).

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Notations

In the following discussion $(\Sigma, \sigma_\Sigma)$ is a connected oriented closed manifold of dimension $2n + 1$. We endow $\Sigma$ with an auxiliary Riemannian metric in order to measure distances and norms of various objects defined on $\Sigma$. We denote by $r$ a real number in $[0, 1]$. We will often consider paths $r \mapsto o_r$ with values in some target space $O$. We use the short-hand $\{o_r\} \subset O$ for such a path and $[o_r]$ for a homotopy class of paths with conveniently chosen boundary conditions. If $\{o_r\}$ and $\{o'_r\}$ are two paths with $o_0 = o'_0$, we write $\{o_r\} \# \{o'_r\}$ for the concatenated path. Finally, if a mathematical object is expressed by a certain symbol, we add a prime to it in order to denote another object of the same kind.

2 The volume of a closed two-form

2.1 Definition of the volume

Let $\Xi^k(\Sigma)$ be the set of closed $k$-forms on $\Sigma$ and $\Xi^2_C(\Sigma)$ the set of elements in $\Xi^2(\Sigma)$ representing the de Rham cohomology class $C \in H^2_{dR}(\Sigma)$. The set $\Xi^2_C(\Sigma)$ is affine with underlying vector space

$$\tilde{\Omega}^1(\Sigma) := \frac{\Omega^1(\Sigma)}{\Xi^1(\Sigma)}.$$ 

Indeed, for every $\Omega \in \Xi^2_C(\Sigma)$, we have the surjective map

$$B_\Omega : \Omega^1(\Sigma) \longrightarrow \Xi^2_C(\Sigma), \quad \alpha \mapsto \Omega + d\alpha.$$ 

This map passes to the quotient under $\Xi^1(\Sigma)$ and yields a bijection

$$\tilde{B}_\Omega : \tilde{\Omega}^1(\Sigma) \longrightarrow \Xi^2_C(\Sigma)$$

so that we have an identification

$$T_\Omega \Xi^2_C(\Sigma) = \tilde{\Omega}^1(\Sigma).$$

We define now an exact one-form $\text{vol} \in \Omega^1(\tilde{\Omega}^1(\Sigma))$. To this purpose, we fix a reference two-form $\Omega_0 \in \Xi^2_C(\Sigma)$, and for every $\alpha \in \Omega^1(\Sigma)$, we use the short-hand

$$\Omega_\alpha := B_{\Omega_0}(\alpha) = \Omega_0 + d\alpha.$$ 

For every $\beta \in \Omega^1(\Sigma) \cong T_\alpha \Omega^1(\Sigma)$, we set

$$\text{vol}_\alpha \cdot \beta := \int_{\Sigma} \beta \wedge \Omega^n_\alpha.$$ 

In the lemma below, we show that the one-form $\text{vol}$ admits the primitive functional

$$\text{Vol} : \Omega^1(\Sigma) \rightarrow \mathbb{R}, \quad \text{Vol}(\alpha) := \int_0^1 \left( \int_{\Sigma} \alpha \wedge \Omega^n_{\tau\alpha} \right) d\tau. \quad (2.1)$$

Remark 2.1. Exchanging the order of integration, we can rewrite

$$\text{Vol}(\alpha) = \int_{\Sigma} \text{CS}(\alpha), \quad \text{CS}(\alpha) := \int_0^1 \alpha \wedge \Omega^n_{\tau\alpha} d\tau.$$
When $C = 0$ and $\Omega_0 = 0$, the form $(n + 1)\text{CS}(\alpha)$ reduces to the Chern-Simons form for principal $S^1$-bundles with invariant polynomial $P(x_1, \ldots, x_{n+1}) = x_1 \cdots x_{n+1}$ \cite{CS74} equation (3.1). Furthermore, $(n + 1)\text{Vol}(\alpha)$ corresponds to the cohomology class in \cite{CS74} Theorem 3.9, case $2l - 1 = n$ (where $M$ therein is our $\Sigma$ and $l$ is our $n + 1$).

**Lemma 2.2.** The functional $\text{Vol}$ associated with $\Omega_0$ is uniquely characterised by the properties

\[ \text{Vol}(0) = 0, \quad d\text{Vol} = \text{vol}. \]

The following formula holds (see also \cite{CS74} equation (3.5)):

\[ \text{Vol}(\alpha) = \sum_{j=0}^{n} \frac{1}{j+1} \binom{n}{j} \int_{\Sigma} \alpha \wedge (d\alpha)^j \wedge \Omega_0^{n-j}. \]

**Proof.** Differentiating under the integral sign we write $d_{\alpha} \text{Vol}(\beta) = \int_0^1 V(r)dr$ and compute

\[ V(r) = \int_{\Sigma} \left( \beta \wedge \Omega_r^\alpha + \alpha \wedge \Omega_r^{n-1} \wedge nr d\beta \right) \]

\[ = \int_{\Sigma} \left( \beta \wedge \Omega_r^\alpha + nr \Omega_r^{n-1} \wedge \left( \beta \wedge d\alpha - d(\alpha \wedge \beta) \right) \right) \]

\[ = \int_{\Sigma} \left( \beta \wedge \Omega_r^\alpha + r \beta \wedge \Omega_r^{n-1} \wedge (n \alpha) \right) \]

\[ = \frac{d}{dr} \left( \int_{\Sigma} r \beta \wedge \Omega_r^\alpha \right), \]

where in the third equality we used Stokes' Theorem. We conclude that

\[ d_{\alpha} \text{Vol} \cdot \beta = \int_0^1 \frac{d}{dr} \left( \int_{\Sigma} r \beta \wedge \Omega_r^\alpha \right) dr = \int_{\Sigma} \frac{1}{2} \beta \wedge \Omega_r^{n-1} - \int_{\Sigma} 0 \beta \wedge \Omega_r^{n-1} = \text{vol}_{\alpha} \cdot \beta. \]

The formula for $\text{Vol}$ follows by expanding the binomial $\Omega_r^\alpha = (\Omega_0 + r d\alpha)^n$ in (2.1) and integrating each term. \hfill \Box

**Remark 2.3.** If dim $\Sigma = 3$, namely $n = 1$, the formula for $\text{Vol}$ turns into

\[ \text{Vol}(\alpha) = \int_{\Sigma} \alpha \wedge (\Omega_0 + \frac{1}{2} d\alpha). \]

Let us specify the dependence of $\text{Vol}$ on the reference form $\Omega_0$.

**Lemma 2.4.** Let $\text{Vol}' : \Omega^1(\Sigma) \to \mathbb{R}$ be the volume functional associated with another reference two-form $\Omega'_0 \in \mathbb{R}^2(\Sigma)$. If $\alpha' \in \Omega^1(\Sigma)$ is such that $\Omega'_0 = \Omega_0 + d\alpha'$, then

\[ \text{Vol}(\alpha) = \text{Vol}(\alpha') + \text{Vol}'(\alpha - \alpha'), \quad \forall \alpha \in \Omega^1(\Sigma). \]

**Proof.** For every $\alpha, \beta \in \Omega^1(\Sigma)$, we have

\[ \text{vol}'_{\alpha - \alpha' \cdot \beta} = \int_{\Sigma} \beta \wedge (\Omega'_0 + d\alpha - d\alpha')^n = \int_{\Sigma} \beta \wedge (\Omega_0 + d\alpha)^n = \text{vol}_{\alpha} \cdot \beta, \]

where $\text{vol}' := d\text{Vol}'$. Therefore, from the fundamental theorem of calculus, we get

\[ \text{Vol}(\alpha) - \text{Vol}(\alpha') = \int_0^1 \text{vol}_{r(\alpha - \alpha') + \alpha'} \cdot (\alpha - \alpha')dr = \int_0^1 \text{vol}'_{(\alpha - \alpha') \cdot (\alpha - \alpha')dr} = \text{Vol}'(\alpha - \alpha'). \]

\hfill \Box
We now wish to use the map \( B_{\Omega_0} \) to push the volume form and functional to the space of closed two-forms in the class \( C \). To this purpose, we observe that \( \text{Vol} \) behaves linearly under the addition of closed one-forms. More precisely, for all \( \alpha \in \Omega^1(\Sigma) \) we have

\[
\text{Vol}(\alpha + \tau) - \text{Vol}(\alpha) = \text{Vol}(\tau) = ([\tau] \cup C^n, [\Sigma]), \quad \forall \tau \in \Xi^1(\Sigma),
\]

(2.2)

and similarly,

\[
\text{vol}_\alpha \cdot \tau = ([\tau] \cup C^n, [\Sigma]), \quad \forall \tau \in \Xi^1(\Sigma).
\]

(2.3)

These formulae suggest to treat two separate cases.

**Case 1:** \( C^n = 0 \).

The volume functional \( \text{Vol} : \Omega^1(\Sigma) \to \mathbb{R} \) passes to the quotient by \( \Xi^1(\Sigma) \) according to (2.2):

\[
\mathcal{V}\text{ol} : \Xi^2_C(\Sigma) \to \mathbb{R}, \quad \mathcal{V}\text{ol}(\Omega_0) = 0.
\]

(2.4)

Following Lemma 2.4 if \( \mathcal{V}\text{ol}' \) denotes the functional obtained from \( \Omega_0' \), then there holds

\[
\mathcal{V}\text{ol}(\Omega) = \mathcal{V}\text{ol}(\Omega_0') + \mathcal{V}\text{ol}'(\Omega), \quad \forall \Omega \in \Xi^2_C(\Sigma).
\]

This means that \( \text{vol} \in \Omega^1(\Omega^1(\Sigma)) \) descends to a one-form

\[
\text{vol} \in \Omega^1(\Xi^2_C(\Sigma)), \quad d\text{vol} = \text{vol},
\]

which is *independent* of the reference two-form \( \Omega_0 \).

**Case 2:** \( C^n \neq 0 \).

The functional \( \text{Vol} \) does not descend to \( \Xi^2_C(\Sigma) \). Indeed, by Poincaré duality, there exists a form \( \tau \in \Xi^1(\Sigma) \) such that

\[
[\tau] \cup C^n \neq 0.
\]

(2.5)

However, we can use the volume function to define a distinguished class of one-forms.

**Definition 2.5.** We say that \( \alpha \in \Omega^1(\Sigma) \) is a **normalised one-form**, if the implication

\[
C^n \neq 0 \implies \text{Vol}(\alpha) = 0
\]

holds true (in particular, when \( C^n = 0 \) as in Case 1, all one-forms are normalised).

The inclusion \( \text{Vol}^{-1}(0) \subset \Omega^1(\Sigma) \) induces a bijection

\[
\frac{\text{Vol}^{-1}(0)}{\sim} \to \Omega^1(\Sigma) \overset{B_{\Omega_0}}{\cong} \Xi^2_C(\Sigma),
\]

(2.6)

where \( \alpha' \sim \alpha'' \) if and only if \( \alpha'' - \alpha' \in \Xi^1(\Sigma) \). Indeed, by (2.2) any class \( \bar{\alpha} := \alpha + \Xi^1(\Sigma) \in \tilde{\Omega}^1(\Sigma) \) has a normalised representative \( \alpha' = \alpha + s\tau \), for some \( s \in \mathbb{R} \). We also observe that, if \( \alpha'' \) is another normalised representative of \( \bar{\alpha} \), then, again by (2.2), we have

\[
[\alpha'' - \alpha'] \cup C^n = 0.
\]

(2.7)
In view of (2.6), we define the volume form and functional on $\Xi^2_C(\Sigma)$ trivially by setting
\[ \mathcal{V}ol : \Xi^2_C(\Sigma) \to \mathbb{R}, \quad \mathcal{V}ol := 0, \quad \text{vol} \in \Omega^1(\Xi^2_C(\Sigma)), \quad \text{vol} := 0. \]

If $\alpha$ is a normalised one-form representing $\bar{\alpha} \in \bar{\Omega}^1(\Sigma)$, then the inclusion $\ker \text{vol}_\alpha \subset \Omega^1(\Sigma)$ induces the surjection $\ker \text{vol}_\alpha \to \bar{\Omega}^1(\Sigma) \cong T_{\alpha}\bar{\Omega}^1(\Sigma)$ due to the existence of $\tau$ satisfying (2.5), and in turn the isomorphism
\[ \ker \text{vol}_\alpha \cong T_{\alpha}\bar{\Omega}^1(\Sigma) \sim T_{\bar{\alpha}}\bar{\Omega}^1(\Sigma), \]
thanks to (2.3) and (2.7).

**Remark 2.6.** For contact forms and Hamiltonian systems the volume functional recovers the following well-known formulae.

- *(Contact forms)* This is an instance of Case 1. Let $\Omega_0 = 0$ and $\alpha \in \Omega^1(\Sigma)$ be a contact form. If we endow $\Sigma$ with the orientation $o_\alpha$, we recover the contact volume up to a constant factor:
\[ \mathcal{V}ol(d\alpha) = \text{Vol}(\alpha) = \frac{1}{n+1} \int_\Sigma \alpha \wedge (d\alpha)^n = \frac{1}{n+1} \text{Volume}(\alpha). \]

- *(Hamiltonian systems)* This is an instance of Case 2. Let $\Sigma = M \times S^1$ and $\Omega_0 = p^*\omega_0$ for some symplectic form $\omega_0$ on $M$, where $p : M \times S^1 \to M$ is the projection on the first factor. Let $\alpha = H d\phi$, where $H : M \times S^1 \to \mathbb{R}$ and $\phi$ is the coordinate on $S^1$. Then, the volume functional recovers the Calabi invariant of $H$
\[ \text{Vol}(H d\phi) = \int_{M \times S^1} (H d\phi) \wedge \omega_0^n = \text{CAL}_{\omega_0}(H) \]
and $\alpha$ is normalised if and only if the Calabi invariant of $H$ vanishes.

### 2.2 The volume is invariant under pull-back and isotopies

In this subsection, we prove two invariance results for the volume. Recall that $\Omega_0$ is the fixed reference form in $\Xi^2_C(\Sigma)$ for $C \in H^2_{dR}(\Sigma)$.

**Proposition 2.7.** Let $\Sigma$ and $\Sigma'$ be two closed oriented manifolds of dimension $2n+1$ and $\Pi : \Sigma' \to \Sigma$ a map of degree $\deg \Pi \in \mathbb{Z}$. Let $C' := \Pi^*C \in H^2_{dR}(\Sigma')$ and $\Omega'_0 := \Pi^*\Omega_0 \in \Xi^2_C(\Sigma')$. If $\text{Vol}$ and $\mathcal{V}ol$ are the volumes associated with $\Omega_0$ and $\text{Vol}'$, $\mathcal{V}ol'$ the volumes associated with $\Omega'_0$, there holds
\[ \text{Vol}' \circ \Pi^* = (\deg \Pi) \cdot \text{Vol}, \quad \mathcal{V}ol' \circ \Pi^* = (\deg \Pi) \cdot \mathcal{V}ol. \]

**Proof.** The statement follows immediately from the definition of the volume and the observation that for all top dimensional forms $\mu$ on $\Sigma$ there holds
\[ \int_{\Sigma'} \Pi^*\mu = \deg \Pi \cdot \int_{\Sigma} \mu. \]

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For the second result, we need a little preparation. Let $\Psi : \Sigma \to \Sigma$ be a diffeomorphism isotopic to the identity $\text{id}_\Sigma$. Since the pull-back $\Psi^*$ of $\Psi$ acts as the identity in cohomology, there is a normalised one-form $\theta \in \Omega^1(\Sigma)$ (determined up to a normalised closed one-form) such that
\[d\theta = \Psi^*\Omega_0 - \Omega_0.\] (2.8)

We define
\[\widehat{\Psi}_\theta^* : \Omega^1(\Sigma) \to \Omega^1(\Sigma), \quad \widehat{\Psi}_\theta^*(\alpha) := \theta + \Psi^*\alpha,
\]
so that the following diagram commutes:
\[
\begin{array}{ccc}
\Omega^1(\Sigma) & \xrightarrow{\widehat{\Psi}_\theta^*} & \Omega^1(\Sigma) \\
\downarrow B_{\Omega_0} & & \downarrow B_{\Omega_0} \\
\Xi^2_C(\Sigma) & \xrightarrow{\Psi^*} & \Xi^2_C(\Sigma)
\end{array}
\] (2.9)

Let $\{\Psi_r\}$ be an isotopy with $\Psi_0 = \text{id}_\Sigma$ and $\Psi_1 = \Psi$. This gives rise to a smooth family of one-forms $\{\theta_r\}$ satisfying
\[\theta_0 = 0, \quad d\theta_r = \Psi_r^*\Omega_0 - \Omega_0, \quad \forall r \in [0, 1].\]

To construct such a family, we just take the time-dependent vector field $X_r$ generating the isotopy $\{\Psi_r\}$ and let $\{\theta_r\}$ be the unique path such that
\[\theta_0 = 0, \quad \dot{\theta}_r = \Psi_r^*(\iota_{X_r}\Omega_0).\]

To ease the notation, if $\alpha \in \Omega^1(\Sigma)$, we use the short-hand
\[\widehat{\Psi}_\theta^*(\alpha) := \widehat{\Psi}_{\theta_r}^*(\alpha) = \theta_r + \Psi_r^*\alpha.\] (2.10)

Every $\theta_r$ is normalised, i.e. $\text{Vol}(\theta_r) = 0$. Indeed, $\text{Vol}(\theta_0) = \text{Vol}(0) = 0$ and
\[\frac{d}{dr}\text{Vol}(\theta_r) = \text{vol}_{\theta_r}(\dot{\theta}_r) = \int_{\Sigma} \Psi_r^*(\iota_{X_r}\Omega_0) \wedge (\Psi_r^*\Omega_0)^n = \int_{\Sigma} (\iota_{X_r}\Omega_0) \wedge \Omega_0^n = 0.
\]
In particular, the form $\tau_1 := \theta - \theta_1$ is closed and normalised and as observed in (2.7), there holds
\[\left[\tau_1\right] \cup C^n = 0
\]
(note that this condition is automatically satisfied if $C^n = 0$).

**Proposition 2.8.** If $\Psi$ is a diffeomorphism isotopic to $\text{id}_\Sigma$ and $\theta$ is any normalised one-form satisfying (2.8), there holds
\[\text{Vol} \circ \widehat{\Psi}_\theta^* = \text{Vol}.
\]
As a consequence, the set of normalised one-forms is $\widehat{\Psi}_\theta^*$-invariant and we have
\[\text{Vol} \circ \Psi^* = \text{Vol}.
\]
Proof. For any $\alpha \in \Omega^1(\Sigma)$, we have
\[
\frac{d}{dr}\hat{\Psi}_r^*(\alpha) = \Psi_r^*(\iota_X, \Omega_0) + \Psi_r^*(\iota_X, \alpha + d(\alpha(X_r))) = \Psi_r^*(\iota_X, \Omega_0 + d(\alpha(X_r))).
\]
Using this relation, we compute
\[
\frac{d}{dr}\text{Vol}(\hat{\Psi}_r^*(\alpha)) = \text{vol}(\hat{\Psi}_r^*(\alpha)) \cdot \frac{d}{dr}\hat{\Psi}_r^*(\alpha) = \int_{\Sigma} \Psi_r^*(\iota_X, \Omega_0 + d(\alpha(X_r))) \wedge (\Psi_r^*\Omega_0)^n
= \int_{\Sigma} \iota_X, \Omega_0 \wedge \Omega_0^n + \int_{\Sigma} d(\alpha(X_r)) \wedge \Omega_0^n
= 0.
\]
We conclude from (2.2) and the relation $[\tau_1] \cup C = 0$ that
\[
\text{Vol}(\hat{\Psi}_0^*(\alpha)) = \text{Vol}(\hat{\Psi}_0^*(\alpha) + \tau_1) = \text{Vol}(\hat{\Psi}_0^*(\alpha)) + \langle [\tau_1] \cup C, [\Sigma] \rangle = \text{Vol}(\hat{\Psi}_0^*(\alpha)) = \text{Vol}(\alpha).
\]
This proves the first identity. The second identity follows from the first one and the commutation relation $B_{\Omega_0} \circ \hat{\Psi}_0 = \Psi^* \circ B_{\Omega_0}$ in (2.9).

3 Odd-symplectic forms

3.1 A couple of definitions

For a closed two-form $\Omega \in \Xi^2(\Sigma)$ on $\Sigma$, we consider the (possibly singular) distribution $\ker \Omega \to \Sigma$ defined by
\[
\ker \Omega := \{(z, u) \in T\Sigma \mid \Omega_z(u, v) = 0, \forall v \in T_z\Sigma\}.
\]
The distribution $\ker \Omega$ is naturally co-oriented by $\Omega^n$ and we orient it combining the orientation on $\Sigma$ with such a co-orientation.

Definition 3.1. We say that $\Omega \in \Xi^2(\Sigma)$ is odd-symplectic if $\ker \Omega \to \Sigma$ is a one-dimensional distribution and denote the set of odd-symplectic forms by $\mathcal{S}(\Sigma) \subset \Xi^2(\Sigma)$. If $C \in H^2_{\text{dR}}(\Sigma)$, we write $\mathcal{S}_C(\Sigma) := \mathcal{S}(\Sigma) \cap \Xi^2_C(\Sigma)$.

Remark 3.2. Any of the following conditions gives an equivalent definition of $\Omega \in \Xi^2(\Sigma)$ being odd-symplectic:

- $\Omega^n$ is nowhere vanishing.
- There exists a one-form $\sigma$ on $\Sigma$ such that $\sigma \wedge \Omega^n$ is a volume form.
- There exists a nowhere vanishing one-form $\sigma$ such that $\Omega|_{\ker \sigma}$ is a non-degenerate bilinear form on $\ker \sigma$.

Definition 3.3. Given a reference form $\Omega_0 \in \Xi^2_C(\Sigma)$ and a pair of one-forms $\alpha_0, \sigma_0 \in \Omega^1(\Sigma)$, we can build an affine map
\[
C^\infty(\Sigma) \to \Xi^2_C(\Sigma),
H \mapsto \Omega_{\alpha_0 + H\sigma_0} = \Omega_0 + d(\alpha_0 + H\sigma_0).
\]
We refer to the image of this map, as the set of $H$-forms (with respect to $\Omega_0$, $\alpha_0$, and $\sigma_0$).
3.2 A stability result

As we will see in Section 7, given $\Omega_0$, $\alpha_0$ and $\sigma_0$ as in Definition (3.3), it is extremely helpful to see, if an odd-symplectic form $\Omega_1 \in \mathcal{S}_C(\Sigma)$ admits a diffeomorphism $\Psi : \Sigma \rightarrow \Sigma$ isotopic to $\text{id}_\Sigma$ such that $\Psi^*\Omega_1$ is an H-form. We provide a criterion for such a diffeomorphism to exist.

**Proposition 3.4.** Let $\{\Omega_r = \Omega_0 + d\alpha_r\}$ be a path in $\mathcal{S}_C(\Sigma)$. Let $\{\sigma_r\}$ be a path in $\Omega^1(\Sigma)$ such that $\sigma_r \wedge \Omega^n_r$ is a volume form on $\Sigma$, which exists according to Remark 3.2. Let us denote by $I_r : (\ker \sigma_r)^* \rightarrow \ker \sigma_r$ the inverse of the map $v \mapsto \iota_v \Omega_{\alpha_r}|_{\ker \sigma_r}$. If

$$
(\dot{\sigma}_r - \iota_{I_r(\sigma_r)}d\sigma_r)|_{\ker \sigma_r} = 0, \hspace{1cm} (3.1)
$$

then there exist an isotopy $\{\Psi_r : \Sigma \rightarrow \Sigma\}$ with $\Psi_0 = \text{id}_\Sigma$ generated by a time-dependent vector field $X_r \in \ker \sigma_r$ and a path of functions $\{H_r : \Sigma \rightarrow \mathbb{R}\}$ such that, for every $r \in [0,1]$,

$$
\dot{\Psi}_r^*\alpha_r = \alpha_0 + H_r\sigma_0 + dK_r, \hspace{1cm} K_r := \int_0^r \alpha_r(X_r) \circ \Psi_r dr'.
$$

Here $\dot{\Psi}_r^*$ is the map in (2.10). In particular, we have

$$
\Psi_r^*\Omega_1 = \Omega_0 + d(\alpha_0 + H_1\sigma_0), \hspace{1cm} \text{Vol}(\alpha_1) = \text{Vol}(\alpha_0 + H_1\sigma_0).
$$

**Proof.** The argument is a refinement of the Gray stability theorem, see e.g. [Gei08, Theorem 2.2.2]. We introduce a vector field $V_r$ on $\Sigma$ uniquely determined by the relations

$$
\iota_{V_r}\Omega_{\alpha_r} = 0, \hspace{1cm} \sigma_r(V_r) = 1.
$$

We define a path of vector fields $r \mapsto X_r$ on $\Sigma$ satisfying the relations

$$
X_r \in \ker \sigma_r, \hspace{1cm} (\iota_{X_r}\Omega_{\alpha_r} + \dot{\sigma}_r)|_{\ker \sigma_r} = 0, \hspace{1cm} (3.2)
$$

which exists and is unique by (i). Applying (ii), we also have

$$
(\iota_{X_r}d\sigma_r + \dot{\sigma}_r)|_{\ker \sigma_r} = 0. \hspace{1cm} (3.3)
$$

Let $r \mapsto \Psi_r$ be the isotopy on $\Sigma$ obtained by integrating $X_r$ and setting $\Psi_0 = \text{id}_\Sigma$. We construct an auxiliary path of functions $r \mapsto J_r$ through

$$
J_0 = 0, \hspace{1cm} \dot{J}_r = (d\sigma_r(X_r, V_r) + \dot{\sigma}_r(V_r)) \circ \Psi_r.
$$

Combining the last equation with (3.3), we arrive at

$$
\iota_{X_r}d\sigma_r + \dot{\sigma}_r = (\dot{J}_r \circ \Psi_r^{-1})\sigma_r.
$$

With this piece of information, we can compute

$$
\frac{d}{dr}(\Psi_r^*\sigma_r) = \Psi_r^*(\iota_{X_r}d\sigma_r + \dot{\sigma}_r) = \Psi_r^*(\dot{J}_r \circ \Psi_r^{-1}\sigma_r) = \dot{J}_r(\Psi_r^*\sigma_r).
$$

Since $r \mapsto e^{\dot{J}_r}\sigma_0$ also satisfies the same ordinary differential equation and coincides with $\Psi_r^*\sigma_r$ at $r = 0$, we see that

$$
\Psi_r^*\sigma_r = e^{J_r}\sigma_0. \hspace{2cm} (3.4)
$$
Next we define the path of functions \( \{ H_r : \Sigma \to \mathbb{R} \} \) by
\[
H_0 = 0, \quad \dot{H}_r = (\dot{\alpha}_r(V_r) \circ \Psi_r)e^{J_r}.
\] (3.5)

Conditions (3.2) and (3.5) together can be rephrased as
\[
\iota_{X_r}\Omega_{\alpha_r} + \dot{\alpha}_r = ((\dot{H}_re^{-J_r}) \circ \Psi_r^{-1})\sigma_r.
\] (3.6)

We are ready to prove the formula for \( \tilde{\Psi}_r^*\alpha_r \) in the statement of the proposition. The identity clearly holds for \( r = 0 \) and we just need to show that the \( r \)-derivatives of both sides coincide for every \( r \):
\[
\frac{d}{dr}(\tilde{\Psi}_r^*\alpha_r) = \Psi_r^*(d(\alpha_r(X_r))\iota_{X_r}d\alpha_r + \dot{\alpha}_r) + \dot{\theta}_r = d(\alpha_r(X_r) \circ \Psi_r) + \Psi_r^*(\iota_{X_r}\Omega_{\alpha_r} + \dot{\alpha}_r)
= dK_r + \dot{H}_re^{-J_r}\Psi_r^*\sigma_r
= d\tilde{K}_r + \dot{H}_r\sigma_r,
\]
where we used that \( \dot{\theta}_r = \Psi_r^*\iota_{X_r}\Omega_0 \) and equations (3.4) and (3.6). Finally, Proposition 2.8 yields the identity \( \text{Vol}(\alpha_1) = \text{Vol}(\alpha_0 + H_1\sigma_0) \).

**Remark 3.5.** The form \( \sigma_r \wedge \Omega^n_r \) is a volume form and condition (3.1) is satisfied in the following two extreme cases:

(a) The form \( \Omega_0 \) vanishes and \( \alpha_r = T_r\sigma_r \) is a contact form for every \( r \), where \( \{ T_r \} \) is some path of real numbers. In this case, \( K_r \) and \( \theta_r \) vanish so that we have the usual Gray stability theorem
\[
\Psi_r^*\alpha_r = (1 + \frac{1}{19}H_r)\alpha_0.
\]

(b) The form \( \sigma_0 \) is closed, there holds \( \sigma_r = \sigma_0 \) for all \( r \), and \( \Omega_r \) non-degenerate on \( \ker \sigma_0 \).

**Remark 3.6.** It would be interesting to find a condition not involving \( r \)-derivatives implying (3.1). Such a condition can indeed be found in cases (a) or (b) in the remark above, where the odd-symplectic forms \( \Omega_r \) are actually stable according to [CM05, Section 2.1].

### 4 Oriented \( S^1 \)-bundles

#### 4.1 \( S^1 \)-bundles and free \( S^1 \)-actions

Let \( \mathcal{P}(\Sigma) \) be the space of all oriented \( S^1 \)-bundles having \( \Sigma \) as total space, up to equivalence. Here we say that two bundles are **equivalent** if they have the same oriented fibres. If \( p : \Sigma \to M \) is an element of \( \mathcal{P}(\Sigma) \), then \( M \) is a closed manifold of dimension \( 2n \). We write \( \Xi^2(M) \) for the space of closed two-forms on \( M \) and \( \Xi^2_c(M) \) for the set of those \( \omega \in \Xi^2(M) \) with \( [\omega] = c \). We orient \( M \) combining the orientation on \( \Sigma \) with the orientation of the \( p \)-fibres. We denote by \( -p \) the bundle obtained from \( p \) by reversing the orientation.

**Definition 4.1.** If \( p \) belongs to \( \mathcal{P}(\Sigma) \), then \( p^{-1}(\text{pt}) \subset \Sigma \) denotes an oriented \( p \)-fibre and \( [p^{-1}(\text{pt})] \in H_1(\Sigma; \mathbb{Z}) \) its integral homology class.
Definition 4.2. A bundle map between oriented $S^1$-bundles $p : \Sigma \to M$ and $p^\vee : \Sigma^\vee \to M^\vee$ is a map $\Pi : \Sigma^\vee \to \Sigma$ diffeomorphically sending every oriented fibre of $p^\vee$ to an oriented fibre of $p$. In this case, we write $p^\vee = \Pi^*p$. A bundle map yields a quotient map $\pi : M^\vee \to M$ between the base manifolds so that $p \circ \Pi = \pi \circ p^\vee$. If $\Pi$ is also a diffeomorphism, we say that $\Pi$ is a bundle isomorphism. When $\Sigma^\vee = \Sigma$, we write $p^\vee = p'$, where $p' : \Sigma \to M'$, and denote by $\Psi : \Sigma \to \Sigma$ a bundle isomorphism between $p$ and $p'$ with quotient map $\psi : M' \to M$. We summarise the properties of $\Pi$ and $\Psi$ in two commutative diagrams.

Let $\mathcal{U}(\Sigma)$ be the space of free $S^1$-actions on $\Sigma$. For $u \in \mathcal{U}(\Sigma)$, we denote by $V = V_u$ the fundamental vector field on $\Sigma$ generated by $u$. We have a natural map

$$\mathcal{U}(\Sigma) \to \mathfrak{P}(\Sigma),$$

which associates to a free $S^1$-action the quotient map onto its orbit space. It is a classical fact that this map is surjective. Indeed, let us fix $p \in \mathfrak{P}(\Sigma)$. Using a partition of unity, we find a vector field $X$ on $\Sigma$, positively tangent to the $p$-fibres at every point. If $T(z) > 0$ is the period of the periodic orbit of $X$ starting at $z \in \Sigma$, then the rescaled vector field $V := T X$ yields the desired $S^1$-action. The set $\mathcal{U}(\Sigma)$ carries a natural $C^1$-topology as a closed subset of the space $\text{Maps}(S^1 \times \Sigma, \Sigma)$ and we endow $\mathfrak{P}(\Sigma)$ with the quotient topology brought by the map (4.2). We write $\mathcal{U}(p)$ for the fibre above $p \in \mathfrak{P}(\Sigma)$. This is a convex space in the following sense. If $V_0$ and $V_1 = T_1 V_0$ are the fundamental vector fields of $u_0, u_1 \in \mathcal{U}(p)$, then

$$V_r := \frac{V_1}{r + (1-r)T_1}, \quad \forall r \in [0,1]$$

is the fundamental vector field of some $u_r \in \mathcal{U}(p)$.

Finally, if $\Lambda(\Sigma)$ denotes the space one-periodic curves in $\Sigma$, we have a map

$$j_p : \Sigma \to \Lambda(\Sigma)$$

associating to a point the $u$-orbit through it for some $u \in \mathcal{U}(p)$. Up to an orientation-preserving change of parametrisation of the elements of $\Lambda(\Sigma)$, this map depends only on $p$.

Definition 4.3. If $p : \Sigma \to M$ belongs to $\mathfrak{P}(\Sigma)$ and $u$ belongs to $\mathcal{U}(p)$, we write

$$e_Z \in H^2(M; \mathbb{Z})$$

for minus the Euler class of $u$, as defined, for example in [Che77]. This class is independent of $u \in \mathcal{U}(p)$, and therefore, we refer to it as minus the Euler class of $p$.

Remark 4.4. The bundle $p$ is trivial (namely admits a global section) if and only if $e_Z = 0$. If $p^\vee : \Sigma^\vee \to M^\vee$ is another bundle with minus Euler class $e_Z^\vee \in H^2(M^\vee; \mathbb{Z})$, then

$$p^\vee = \Pi^*p \implies e_Z^\vee = \pi^*e_Z,$$

where $\Pi$ is a bundle map as in (4.1).
The inclusion map $\mathbb{Z} \hookrightarrow \mathbb{R}$ induces a map on the level of homology and cohomology and we write $e$ and $[p^{-1}(pt)]$ for the images of $e_\mathbb{Z}$ and $[p^{-1}(pt)]_\mathbb{Z}$, respectively:

$$
H^2(M; \mathbb{Z}) \rightarrow H^2(M; \mathbb{R}) \cong H^2_{dR}(M), \quad \quad \quad H_1(\Sigma; \mathbb{Z}) \rightarrow H_1(\Sigma; \mathbb{R}).
$$

$$
e_\mathbb{Z} \mapsto e, \quad \quad \quad [p^{-1}(pt)]_\mathbb{Z} \mapsto [p^{-1}(pt)].
$$

From the piece of the Gysin exact sequence $H^0_{dR}(M) \xrightarrow{\omega_\mathbb{Z}} H^2_{dR}(M) \xrightarrow{p^*} H^2_{dR}(\Sigma)$, we know that

$$
\mathbb{R} \cdot e = \ker \left( p^*: H^2_{dR}(M) \rightarrow H^2_{dR}(\Sigma) \right).
$$

Moreover, by the universal coefficient theorem, we get

$$
e_\mathbb{Z} \text{ is torsion } \iff e = 0, \quad \quad \quad [p^{-1}(pt)]_\mathbb{Z} \text{ is torsion } \iff [p^{-1}(pt)] = 0.
$$

The Chern-Weil theory yields a convenient description of the real cohomology class $e$. Let $K(u)$ denote the space of connection one-forms associated with $u \in \Omega(p)$:

$$
K(u) := \{ \eta \in \Omega^1(\Sigma) \mid \eta(V) = 1, \quad d\eta = p^*\kappa_\eta \text{ for some } \kappa_\eta \in \Omega^2(M) \}.
$$

This space is non-empty, and for every $\eta \in K(u)$, the form $\kappa_\eta$ is closed and we have

$$
e = [\kappa_\eta] \in H^2_{dR}(M).
$$

Conversely, for any $\kappa \in \Omega^2(M)$ with $e = [\kappa]$, there is $\eta_\kappa \in K(u)$ with $\kappa_\eta = \kappa$. We denote by

$$
K(p) := \bigcup_{u \in \Omega(p)} K(u)
$$

the set of connection one-forms for $p$.

### 4.2 Flat $S^1$-bundles and the local structure of $\mathfrak{P}(\Sigma)$

This subsection is devoted to the proof of three lemmas. In the first two, we study flat bundles, i.e. with $e = 0$. In the last one, we prove a theorem of Weinstein [Wei74] showing that the space $\mathfrak{P}(\Sigma)$ is locally trivial.

**Lemma 4.5.** Let $p : \Sigma \rightarrow M$ be an oriented $S^1$-bundle. We have an equivalence

$$
e = 0 \iff [p^{-1}(pt)] \neq 0.
$$

(4.5)

If $c \in H^2_{dR}(M)$ and we define $C := p^*c \in H^2_{dR}(\Sigma)$, there holds

$$
\text{PD}(C^n) = \langle c^n, [M] \rangle \cdot [p^{-1}(pt)] \in H_1(\Sigma; \mathbb{R}).
$$

(4.6)

In particular, if $c^n \neq 0$, then

$$
\begin{align*}
(i) \quad & \ker \left( H^1_{dR}(\Sigma) \xrightarrow{\cdot [p^{-1}(pt)]} \mathbb{R} \right) = \ker \left( H^1_{dR}(\Sigma) \xrightarrow{\cdot c^n} H^2_{dR}(\Sigma) \right), \\
(ii) \quad & C^n \neq 0 \iff e = 0.
\end{align*}
$$

(4.7)
Proof. The proof of (4.5) relies on the Gysin sequence and of its Poincaré dual

\[
\begin{array}{c}
H^{2n-2}_{\text{dR}}(M) \xrightarrow{\cap e} H^{2n}_{\text{dR}}(M) \xrightarrow{p^*} H^{2n}_{\text{dR}}(\Sigma) \xrightarrow{p_*} H^{2n-1}(M),
\end{array}
\]

where PD denotes Poincaré duality and \( p_* \) stands for the integration along fibres. We claim that PD\((p^*)\) sends the class of a point \([pt]\) to \([p^{-1}(pt)]\). To this purpose, let \( \mu \in \Xi^{2n}(M) \) be such that PD\(([\mu]) = [pt]\), so that \( \int M \mu = 1 \). Going around the central square in the diagram above, the claim follows, if we can show that PD\((p^*[\mu]) = [p^{-1}(pt)]\). This last equality is true since for all \( \tau \in \Xi^1(\Sigma) \) we have

\[
\langle [\tau], [p^{-1}(pt)] \rangle = p_*[\tau] = p_*(\tau) \cdot \mu = \int_M p_*[\tau] \cdot \mu = \int_{\Sigma} \tau \wedge p^* \mu = \langle [\tau] \cup p^*[\mu], [\Sigma] \rangle,
\]

where in \((\ast)\) we have used Fubini’s Theorem. Therefore, \([p^{-1}(pt)] \neq 0\) if and only if the map \( \cap e \) is equal to zero. This happens if and only if \( \langle e, H_2(M; \mathbb{R}) \rangle = 0 \) as one sees identifying \( H_0(M; \mathbb{R}) \) with \( \mathbb{R} \). Finally, from the universal coefficient theorem, \( \langle e, H_2(M; \mathbb{R}) \rangle = 0 \) if and only if \( e = 0 \).

Let us show (4.6). If \( \omega \in \Xi^2_+(M) \) and \( \tau \in \Xi^1(\Sigma) \) is arbitrary, we have

\[
\langle c^n, [M] \rangle \cdot \langle [\tau], [p^{-1}(pt)] \rangle = \int_M p_*[\tau] \cdot \omega^n = \int_{\Sigma} \tau \wedge p^* \omega^n = \langle [\tau] \cup C^n, [\Sigma] \rangle.
\]

Let us assume that \( \langle c^n, [M] \rangle \neq 0 \). Looking again at the last equation, we readily see that (i) in (4.7) holds. The equivalence in (ii) stems from a combination of (4.6) and (4.5)

\[
C^n \neq 0 \iff \text{PD}(C^n) \neq 0 \iff [p^{-1}(pt)] \neq 0 \iff e = 0. \tag{4.6}
\]

We now show that a flat bundle can always be pulled back to a trivial one.

Lemma 4.6. Let \( p : \Sigma \to M \) be an oriented \( S^1 \)-bundle. The class \( e \in H^2_{\text{dR}}(M) \) vanishes if and only if there exists a finite cyclic cover \( \pi : M^\vee \to M \) such that the pull-back bundle \( p^\vee : \Sigma^\vee \to M^\vee \) of \( p \) by \( \pi \) is trivial. In this case, the order of \( e_Z \) in \( H^2(M; \mathbb{Z}) \) equals the minimum degree of a cover with the properties above.

Proof. We preliminarily observe that if \( \pi : M^\vee \to M \) is a finite cover of degree \( k \) and the bundle \( p^\vee : \Sigma^\vee \to M^\vee \) is the pull-back of \( p \) through \( \pi \) with minus the Euler class \( e_Z^\vee \), then by Remark 4.3 there holds

\[
\pi_* e_Z^\vee = k \cdot e_Z,
\]

where \( \pi_* : H^2(M^\vee; \mathbb{Z}) \to H^2(M; \mathbb{Z}) \) is the transfer map (see [Hat02, Chapter 3.G]). Therefore, if we suppose that \( p^\vee \) is trivial, we deduce that \( e_Z^\vee = 0 \), and hence, \( k \cdot e_Z = 0 \). In particular, the order of \( e_Z \) divides \( k \).

Conversely, let \( k \) be a positive integer such that \( k \cdot e_Z = 0 \) and take \( u \in \mu_k \). Let \( j : \Sigma \to S^1 \) be the canonical homomorphism \( j : j \to j/k \), and we denote by \( u_k \) the free \( \mathbb{Z}/k \)-action on \( \Sigma \) obtained combining this homomorphism with \( u \). Let \( \Sigma_k := \Sigma/\mu_k \) be the quotient by this action and denote by \( \Pi_k : \Sigma \to \Sigma_k \) the associated quotient map. The natural map \( p_k : \Sigma_k \to M \) is an oriented \( S^1 \)-bundle. Let \( (e_Z)_k \in H^2(M; \mathbb{Z}) \) denote minus the Euler class of
that leaves $M^\lor$ invariant. The deck transformation group is isomorphic to $\mathbb{Z}/k\mathbb{Z}$, where $j$ divides $k$, since it is a subgroup of the cyclic group $\mathbb{Z}/k\mathbb{Z}$. Let $\Pi : \Sigma^\lor \to \Sigma$ be a covering map lifting $\pi$. We sum up the construction above in the commutative diagram

$$
\begin{array}{c}
\Sigma^\lor \ar@{->}[r]^\Pi \ar@{->}[d]_s & \Sigma \ar@{->}[d]_{p_k} \\
M^\lor \ar@{->}[r]^\pi & M
\end{array}
$$

We construct a section $s^\lor : M^\lor \to \Sigma^\lor$ tautologically, as follows. Let us consider an arbitrary $q^\lor \in M^\lor$. By definition of pull-back bundle, the $p^\lor$-fibre over $q^\lor$ is identified through $\Pi$ with the $p$-fibre over $\pi(q^\lor)$. Since $i(q^\lor) \in \Sigma$ belongs to the $p$-fibre over $\pi(q^\lor)$, we can define $s^\lor(q^\lor)$ through the equation

$$
\Pi(s^\lor(q^\lor)) = i(q^\lor).
$$

This proves the existence of a cover as in the statement of the lemma and shows that the degree $k^\lor$ is less than or equal to the order of $e_Z$. 

If $p \in \Psi(\Sigma)$ and $\{\Psi_r : \Sigma \to \Sigma\}$ is an isotopy with $\Psi_0 = \text{id}_\Sigma$, then $\{\Psi_r := \Psi_r^* p\}$ yields a path in $\Psi(\Sigma)$. The next lemma shows that all paths in $\Psi(\Sigma)$ arise in this way.

**Lemma 4.7.** The following two statements hold:

(i) Every oriented $S^1$-bundle $p_0 \in \Psi(\Sigma)$ has a $C^1$-neighbourhood $W$ and a continuous map $(p \in W) \mapsto (\Psi_p \in \text{Diff}(\Sigma))$ such that

- $\Psi_{p_0} = \text{id}_\Sigma$,
- $\Psi_p^* p = p_0, \forall p \in W$.

(ii) If $\{p_r : \Sigma \to M_r\}$ is a path in $\Psi(\Sigma)$, there exists an isotopy $\{\Psi_r\}$ of $\Sigma$ such that

- $\Psi_0 = \text{id}_\Sigma$,
- $\Psi_r^* p_r = p_0, \forall r \in [0,1]$.

**Proof.** Pick an arbitrary connection $\eta \in K(p_0)$ and an arbitrary Riemannian metric on $M$ with injectivity radius $\rho_{\text{inj}}$ and distance function $\text{dist} : M \times M \to [0,\infty)$. We define the open subset of $M \times \Sigma$

$$
W := \left\{(q,z) \in M \times \Sigma \bigg| \text{dist}(q,p_0(z)) < \rho_{\text{inj}}/2\right\}
$$

and let $q_M : W \to M$ and $\Sigma : W \to \Sigma$ be the projections on the two factors. We define $W$ as the space of all oriented $S^1$-bundles $p$ such that the following two properties hold:

- For every $z_1, z_2$ in $\Sigma$ belonging to the same $p$-fibre, we have $\text{dist}(p_0(z_1), p_0(z_2)) < \rho_{\text{inj}}/2$;
- For every $z \in \Sigma$, the disc $D_z \subset \Sigma$ is transverse to the orbits of $p$. Here, $D_z$ is the union of all the $\eta$-horizontal lifts through $z$ of the geodesic rays in $M$ emanating from $p_0(z)$ and with length $\rho_{\text{inj}}$. 

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For every \( p \in \mathcal{W} \), we construct a map of fibre bundles over \( M \) given by

\[
S : W \to TM, \quad S(q, z) := (q, v(q, z)), \quad v(q, z) := \int_{S^1} \exp_q^{-1} \circ p_0 \circ \gamma_z(t) dt,
\]

where \( \gamma_z : S^1 \to \Sigma \) is the oriented \( p \)-fibre passing through \( z \). By definition, \( S \) is constant along the \( p \)-fibres. Let \( 0_M \subset TM \) be the zero section and set \( G := S^{-1}(0_M) \). We claim that

- \( S \) is transverse to \( 0_M \),
- \( p_M := q_M|_G : G \to M \) is an \( S^1 \)-bundle map,
- \( \Psi_\Sigma := 3_\Sigma|_G : G \to \Sigma \) is a diffeomorphism.

This is clear if \( p = p_0 \), since in this case \( S(q, z) = \exp_q^{-1}(p_0(z)) \) and \( G = \{(q, z) \mid p_0(z) = q\} \). Therefore, up to shrinking \( \mathcal{W} \), it also holds for \( p \in \mathcal{W} \), as \( S \) depends continuously on \( p \) and being transverse, or a submersion or a diffeomorphism is a \( C^1 \)-open condition \cite[Chapter 2]{Hir94}. Thus, we get an \( S^1 \)-bundle \( p' := p_M \circ \Psi_\Sigma^{-1} : \Sigma \to M \), which is equivalent to \( p \). We now define a map \( \Phi_p : \Sigma \to \Sigma \) such that \( \Phi_p^* p_0 = p \). For every \( z \in \Sigma \), \( \Phi_p(z) \) is the parallel transport with respect to \( \eta \) along the radial geodesic on \( M \) connecting \( p_0(z) \) with \( p'(z) \). The dependence of \( \Phi_p \) from \( p \) is continuous and clearly \( \Phi_p^* p_0 = \text{id}_\Sigma \). As being a diffeomorphism is a \( C^1 \)-open condition, we have \( \Phi_p \in \text{Diff}(\Sigma) \) and (i) is proved by setting \( \Psi_p = \Phi_p^{-1} \).

For part (ii), we break the given path into short paths and apply the first part. \( \square \)

5 Weakly Zoll pairs and Zoll odd-symplectic forms

5.1 Definitions and first properties

Definition 5.1. A couple \((p, c)\) is called a weakly Zoll pair if \( p : \Sigma \to M \) is an element of \( \mathcal{F}(\Sigma) \) and \( c \in H^2_{\text{dR}}(M) \). We say that \( \Omega \in \Xi^2(\Sigma) \) is associated with \((p, c)\), if \( \Omega = p^* \omega \) for some \( \omega \in \Xi^2_c(M) \). We denote by \( 3(\Sigma) \) the set of weakly Zoll pairs and by \( 3_C(\Sigma) \) the subset of those \((p, c) \in 3(\Sigma) \) with \( p^* c = C \in H^2_{\text{dR}}(\Sigma) \).

Let \((p_0, c_0) \in 3_C(\Sigma) \) with \( p_0 : \Sigma \to M_0 \) and take any \( \omega_0 \in \Xi^2_c(M_0) \). We consider the volume functionals \( \text{Vol} : \Omega^1(\Sigma) \to \mathbb{R} \) and \( \mathcal{V} \text{ol} : \Xi^2_C(\Sigma) \to \mathbb{R} \) defined in Section \ref{sec:2} with respect to \( \Omega_0 := p_0^* \omega_0 \in \Xi^2_c(\Sigma) \). We use \( \mathcal{V} \text{ol} \) to get a volume on \( 3_C(\Sigma) \) denoted with the same name:

\[
\mathcal{V} \text{ol} : 3_C(\Sigma) \to \mathbb{R}, \quad \mathcal{V} \text{ol}(p, c) := \mathcal{V} \text{ol}(p^* \omega), \quad \omega \in \Xi^2_c(M).
\]

Lemma 5.2. The following three statements hold.

(i) The volume functional \( \mathcal{V} \text{ol} : 3_C(\Sigma) \to \mathbb{R} \) is well-defined.

(ii) For all \( \zeta \in \Omega^1(M_0) \), \( p_0^* \zeta \) is \( p_0^* \omega_0 \)-normalised, namely

\[
\text{Vol}(p_0^* \zeta) = 0.
\]

(iii) The functions \( \mathcal{V} \text{ol} : \Xi^2_C(\Sigma) \to \mathbb{R}, \) \( \mathcal{V} \text{ol} : 3_C(\Sigma) \to \mathbb{R} \) depend only on the pair \((p_0, c_0)\) and not on the chosen \( \omega_0 \in \Xi^2_{c_0}(M_0) \).
Proof. All three items will stem out a preliminary result. Let \((p, c) \in \mathfrak{Z}_C(\Sigma)\) with \(\omega \in \Xi_2^2(M)\) and pick \(\alpha_\omega \in \Omega^1(\Sigma)\) such that \(p_0^*\omega_0 + d\alpha_\omega = p^*\omega\). We claim that
\[
\text{Vol}(\alpha_\omega + p^*\zeta) = \text{Vol}(\alpha_\omega), \quad \forall \zeta \in \Omega^1(M).
\] (5.1)
Indeed, if we set \(\alpha_r := \alpha_\omega + rp^*\zeta\) so that \(p_0^*\omega_0 + d\alpha_r = p^*(\omega + rd\zeta)\) and \(\dot{\alpha}_r = p^*\zeta\), then
\[
\text{Vol}(\alpha_\omega + p^*\zeta) - \text{Vol}(\alpha_\omega) = \int_0^1 \text{vol}_{\alpha_r} \cdot \dot{\alpha}_r \, dr = \int_0^1 \left( \int_{\Sigma} p^* \zeta \wedge p^*(\omega + rd\zeta)^n \right) \, dr = 0.
\]
If \(\omega'\) is another form in \(\Xi_2^2(M)\), then there is \(\zeta \in \Omega^1(M)\) with \(\omega' - \omega = d\zeta\) and (5.1) implies
\[
\mathfrak{Vol}(p^*\omega') = \mathfrak{Vol}(p^*\omega),
\]
which establishes item (i). Choosing \((p, c) = (p_0, c_0)\) and \(\alpha = 0\) in (5.1), we get item (ii). For item (iii), we take another form \(\omega'_0 \in \Xi_2^2(M_0)\) and write \(\omega'_0 - \omega = d\zeta_0\) for some \(\zeta_0 \in \Omega^1(M_0)\). Applying Lemma 2.4 with \(\Omega_0 = p_0^*\omega_0\) and \(\Omega'_0 = p_0^*\omega'_0\) together with item (ii), we deduce
\[
\text{Vol}(\alpha) = \text{Vol}'(\alpha - p_0^*\zeta_0), \quad \forall \alpha \in \Omega^1(\Sigma),
\]
where \(\text{Vol}'\) is the volume functional associated with \(\Omega'_0\). This implies \(\mathfrak{Vol}(\Omega) = \mathfrak{Vol}'(\Omega)\) for all \(\Omega \in \Xi_2^2(\Sigma)\). \(\square\)

We have a canonical projection
\[
\Psi : \mathfrak{Z}(\Sigma) \to \mathfrak{Z}(\Sigma), \quad \Psi(p, c) = p.
\]
Let us fix a class \(C \in H_{\text{dR}}^2(\Sigma)\) and a connected component \(\mathfrak{Z}^0(\Sigma)\) of \(\mathfrak{Z}(\Sigma)\). We define
\[
\mathfrak{Z}^0_C(\Sigma) := \Psi^{-1}(\mathfrak{Z}^0(\Sigma)) \cap \mathfrak{Z}_C(\Sigma).
\]
We consider the restriction of \(\Psi\) on this set
\[
\Psi^0_C : \mathfrak{Z}^0_C(\Sigma) \to \mathfrak{Z}^0(\Sigma).
\]
By (1.1), for every \((p, c) \in \mathfrak{Z}^0_C(\Sigma)\), we have a surjective map
\[
\mathbb{R} \to (\Psi^0_C)^{-1}(p), \quad A \mapsto (p, Ae + c),
\]
where \(e\) is minus the real Euler class of \(p\). Finally, we define the evaluation map
\[
\text{ev} : \mathfrak{Z}^0_C(\Sigma) \to \mathbb{R}, \quad \text{ev}(p, c) := \langle e^n, [M] \rangle.
\]

Definition 5.3. We say that \(\mathfrak{Z}^0_C(\Sigma)\) is non-degenerate, if the map \(\text{ev} : \mathfrak{Z}^0_C(\Sigma) \to \mathbb{R}\) is non-zero, namely if there exists \((p, c) \in \mathfrak{Z}^0_C(\Sigma)\) with \(\text{ev}(p, c) \neq 0\).

Lemma 5.4. Let \(p_0 : \Sigma \to M_0\) be an element in \(\mathfrak{Z}^0(\Sigma)\), and let \(e_0 \in H_{\text{dR}}^2(M_0)\) be minus the real Euler class of \(p_0\). For every \(p : \Sigma \to M\) in \(\mathfrak{Z}^0(\Sigma)\), there exists a diffeomorphism \(\Psi : \Sigma \to \Sigma\) isotopic to \(\text{id}_\Sigma\) such that \(\Psi^*p = p_0\). If \(\psi : M_0 \to M\) is the quotient map of \(\Psi\) (see (1.1)), then, for every \((p, c) \in (\Psi^0_C)^{-1}(p)\) and \(A \in \mathbb{R}\), there holds
\[
\begin{align*}
(i) \quad & \psi^*(Ae + c) = Ae_0 + \psi^*c, \\
(ii) \quad & \text{ev}(p, Ae + c) = \text{ev}(p_0, Ae_0 + \psi^*c), \\
(iii) \quad & \mathfrak{Vol}(p, Ae + c) = \mathfrak{Vol}(p_0, Ae_0 + \psi^*c).
\end{align*}
\]
Proof. Since $\mathcal{P}^{0}(\Sigma)$ is connected, the existence of a diffeomorphism $\Psi$ as in the statement follows from Lemma [4.7](ii). By Re, we see that $\psi^* c = e_0$, which immediately implies item (i). For item (ii), we observe that $[M] = \psi_s [M_0]$ since $\psi$ is an orientation-preserving diffeomorphism, and compute

$$\text{ev}(p, Ae + c) = \langle (Ae + c)^n, \psi_* [M_0] \rangle = \langle \psi^* (Ae + c)^n, [M_0] \rangle = \text{ev}(p_0, Ae_0 + \psi^* c).$$

For the last relation, we take any $\omega_A \in Z^2(M)$ such that $[\omega_A] = Ae + c$. Then $\Omega_A := p^* \omega_A$ is associated with $(p, Ae + c)$. The form $\Psi^* \Omega_A \in Z^2_C(\Sigma)$ is associated with $(p_0, Ae_0 + \psi^* c)$. Indeed, $\Psi^* \Omega_A = \Psi^* p^* \omega_A = p_0^* \psi^* \omega_A$ and $[\psi^* \omega_A] = \psi^* (Ae + c) = Ae_0 + \psi^* c$. Thus, from Proposition [2.8] we derive

$$\mathcal{V}ol(p, Ae + c) = \mathcal{V}ol(\Omega_A) = \mathcal{V}ol(\Psi^* \Omega_A) = \mathcal{V}ol(p_0, Ae_0 + \psi^* c). \qed$$

Corollary 5.5. If $\mathcal{Z}_C^0(\Sigma)$ is non-empty, the four statements below hold.

(i) The real Euler class of some bundle in $\mathcal{Z}_C^0(\Sigma)$ vanishes, if and only if the real Euler class of every element of $\mathcal{Z}_C^0(\Sigma)$ vanishes.

(ii) The set $\mathcal{Z}_C^0(\Sigma)$ is non-degenerate, if and only if, for every $p \in \mathcal{Z}_C^0(\Sigma)$ there exists a pair $(p, c) \in \mathcal{Z}_C^0(\Sigma)$ such that $\text{ev}(p, c) \neq 0$.

(iii) If the real Euler class of the bundles in $\mathcal{Z}_C^0(\Sigma)$ vanishes, then $\mathcal{Z}_C^0$ is a diffeomorphism and $\text{ev} : \mathcal{Z}_C^0(\Sigma) \to \mathbb{R}$ is a constant map.

(iv) If the real Euler class of the bundles in $\mathcal{Z}_C^0(\Sigma)$ does not vanish, then $\mathcal{Z}_C^0$ has the structure of an affine $\mathbb{R}$-bundle. The $\mathbb{R}$-action on some $(p, c) \in \mathcal{Z}_C^0(\Sigma)$ is given by

$$A \cdot (p, c) = (p, Ae + c), \quad \forall A \in \mathbb{R}.$$

Proof. Items (i) and (ii) are direct consequences of items (i) and (ii) in Lemma [5.1].

To prove item (iii), let us assume that bundles in $\mathcal{Z}_C^0(\Sigma)$ have vanishing real Euler class. By (4.4), the $\mathcal{P}_C^0$-fibres are sets containing only one element. This shows that $\mathcal{P}_C^0$ is a diffeomorphism. This together with item (ii) in Lemma [5.1] yields that $\text{ev}$ is constant.

We suppose that the real Euler class does not vanish and prove item (iv). By (4.4), the $\mathcal{P}_C^0$-fibres are in bijection with $\mathbb{R}$ through the maps $A \mapsto (p, Ae + c)$. We construct an explicit local trivialisation of the bundle structure. Lemma [4.7](i) yields a neighbourhood $W$ of $p_0 \in \mathcal{P}_C^0(\Sigma)$ and a map $p \mapsto \Psi_p$ from $W$ to a neighbourhood of $id_\Sigma$ inside $\text{Diff}(\Sigma)$ such that $\Psi_{p_0} = id_\Sigma$ and $\Psi_p p = p_0$. We define $\Phi := \Psi^{-1}$ and let $\phi_p : M \to M_0$ be the quotient map. If $(p_0, c_0) \in (\mathcal{P}_C^0)^{-1}(p_0)$, then the map

$$W \times \mathbb{R} \to (\mathcal{P}_C^0)^{-1}(W), \quad (p, A) = (p, Ae + \phi_p c_0)$$

provides a local trivialisation around $p_0$. \qed

Definition 5.6. A two-form $\Omega$ on $\Sigma$ is called Zoll, if it is odd-symplectic and there exists $p_0 : \Sigma \to M_0$ in $\mathcal{B}(\Sigma)$ such that the oriented $p_0$-fibres are positively tangent to $ker \Omega$. In this case, we say that $p_0$, which is determined up to equivalence, is the bundle associated with $\Omega$. We write $Z(\Sigma)$ for the set of Zoll (odd-symplectic) forms on $\Sigma$ and $Z_C(\Sigma)$ for the subset of those forms with class $C \in H^2_{\text{dR}}(\Sigma)$. 26
If \( \Omega \) is Zoll, then the form \( \Omega \) descends to the base manifold \( M_\Omega \), since \( d\Omega = 0 \). Namely, there exists \( \omega_\Omega \in \Omega^2(M_\Omega) \) such that

\[
\Omega = p_\Omega^* \omega_\Omega.
\]

The form \( \omega_\Omega \) is closed as well, since \( p_\Omega^* \) is injective. Furthermore, \( \omega_\Omega \) is symplectic since \( \Omega \) is odd-symplectic and the orientations induced by \( \omega_\Omega \) and by \( \omega_\Omega \) on \( M_\Omega \) coincide. We have a natural inclusion

\[
\mathcal{Z}(\Sigma) \to \mathfrak{z}(\Sigma), \quad \Omega \mapsto (p_\Omega, [\omega_\Omega]).
\]

As \( \omega_\Omega \) is positive symplectic, we deduce that

\[
\text{ev}(p_\Omega, [\omega_\Omega]) > 0.
\]

**Remark 5.7.** The inequality above implies that for \( \Omega \in \mathcal{Z}(\Sigma) \) the component \( \mathcal{Z}^0(\Sigma) \) which \((\pm p_\Omega, [\omega_\Omega])\) is belonging to, is non-degenerate.

**Remark 5.8.** By [BW58], Zoll forms with vanishing cohomology class are just the exterior differentials of Zoll contact forms. Indeed, if \( \Omega = p_\Omega^* \omega_\Omega \) is exact, then the cohomology class of \( \omega_\Omega \) is non-zero, as \( \omega_\Omega \) is symplectic, and lies in the kernel of the map \( p_\Omega^* \). In view of (4.4), this means that there exists \( T \in \mathbb{R} \setminus \{0\} \) such that \( \frac{d}{dT} \omega_\Omega = e \). Then, there exists a connection \( \eta \in \mathcal{K}(p_\Omega) \) with \( \kappa_\eta = \frac{1}{T} \omega_\Omega \) and \( \alpha := T \eta \) is a Zoll contact form with \( d\alpha = \Omega \). The orientations \( \sigma_\Sigma \) and \( \sigma_\alpha \) coincide exactly when \( T > 0 \).

Before moving further, it is worthwhile to briefly discuss stability properties of Zoll forms. Let \( \{\Omega_r\} \) be a path in \( \mathcal{Z}_C(\Sigma) \) with the corresponding path of associated bundles \( \{p_r\} \). If the real Euler class of the bundles vanishes, we aim at finding an isotopy \( \{\Psi_r\} \) of \( \Sigma \), such that

\[
\Psi_r^* \Omega_r = \Omega_0.
\]

If the real Euler class of the bundles does not vanish, we aim at finding an isotopy \( \{\Psi_r\} \) of \( \Sigma \), a path of real numbers \( \{A_r\} \) with \( A_0 = 0 \), and \( \eta_0 \in \mathcal{K}(p_0) \) such that

\[
\Psi_r^* \Omega_r = \Omega_0 + A_r d\eta_0.
\]

In this last case, it seems unlikely that all paths admit such an expression. That this happens if \( C = 0 \) is a result of Weinstein [Wei74]. The stability for \( e = 0 \) is reminiscent of [Gin87].

**Proposition 5.9.** Let \( \{\Omega_r\} \) be a path in \( \mathcal{Z}_C(\Sigma) \) with the path of associated bundles \( \{p_r\} \).

- If \( C = 0 \), there is \( \eta_0 \in \mathcal{K}(p_0) \), an isotopy \( \{\Psi_r\} \) of \( \Sigma \) and real numbers \( \{T_r\} \) such that
  \[
  \Psi_r^* \Omega_r = T_r d\eta_0 = \Omega_0 + (T_r - T_0)d\eta_0.
  \]

- If \( e = 0 \), there is an isotopy \( \{\Psi_r\} \) of \( \Sigma \) such that \( \Psi_r^* \Omega_r = \Omega_0 \).

**Proof.** By Lemma 4.7, we can suppose in both cases that \( \{\Omega_r = p_r^* \omega_r\} \), where \( \omega_r \) is a path of symplectic forms on \( M \). If \( C = 0 \), [KN96] yield \( \{\eta_r\} \subset \mathcal{K}(p_0) \) and non-zero real numbers \( \{T_r\} \) such that \( \Omega_r = d(T_r \eta_r) \) and \( \eta_r = \eta_0 + p_r^* \zeta_r \), where \( \{\zeta_r\} \) are one-forms on \( M \). Setting \( \alpha_r := T_r \eta_r \) and applying the Gray stability theorem from Remark 5.5(a), we see that there exists an isotopy \( \{\Psi_r\} \) and a path \( \{H_r : \Sigma \to \mathbb{R}\} \) with \( H_0 = 0 \) such that

\[
\Psi_r^* \alpha_r = (1 + \frac{1}{T_0} H_r) \alpha_0.
\]

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Since \( \dot{\alpha}_r = \dot{T}_r \eta_r + T_r p^* \dot{\zeta}_r \), equation (3.5) implies \( \dot{H}_r = \dot{T}_r \) and hence \( H_r = T_r - T_0 \), which readily implies the desired formula.

If \( e = 0 \), then \( \omega_r = \omega_0 + d \zeta_r \) due to (4.4). By Remark 3.5(b) and Proposition 3.4 with \( \alpha_0 = 0 \), it follows that \( \Psi^* \Omega_r = \Omega_0 + d(H_r \eta_0) \), for some paths \( \{ \Psi_r \} \) and \( \{ H_r \} \) as above. Again by equation (3.5), we conclude that \( H_r = 0 \) for all \( r \).

We are now ready to classify Zoll odd-symplectic forms on a three-dimensional manifold, as promised in the introduction.

5.2 Classification of Zoll odd-symplectic forms in dimension three

Proof of Proposition 1.9

Let \( \Sigma \) have dimension three, let \( b_\Sigma \) be the rank of the free part of \( H_1(\Sigma; \mathbb{Z}) \), and let \( |H_1^{\text{tor}}(\Sigma; \mathbb{Z})| \) denote the cardinality of its torsion subgroup. Let \( \Omega \in Z_C(\Sigma) \) be a Zoll form with cohomology class \( C \) and let \( p = p_\Omega \) its associated bundle. If \( C = 0 \), we know from Remark 5.8 that \( \Omega \) is the differential of a Zoll contact form. In particular, \( \Sigma \) is the total space of a non-trivial oriented \( S^1 \)-bundle, \( b_\Sigma \) is even and we already treated this case in [BK19a, Proposition 1.2].

Suppose that \( C \neq 0 \). In this case, \( e = 0 \) by equivalence (ii) in (4.7). Since \( M \) is a surface this implies that \( e_\Sigma = 0 \), and hence \( p \) is trivial. Therefore, \( \Sigma \cong M \times S^1 \) and we see that

\[
b_\Sigma = 1 + 2 \text{ genus}(M), \quad |H_1^{\text{tor}}(\Sigma; \mathbb{Z})| = 1. \tag{5.2}\]

In particular, \( b_\Sigma \) is odd. Let \( \Omega' \in Z_C(\Sigma) \) be another Zoll form with class \( C' \). Since \( b_\Sigma \) is odd, then \( C' \neq 0 \) and \( p' : \Sigma \to M' \) is the trivial bundle. We write \( \Omega = p^* \omega \) and \( \Omega' = p'^* \omega' \), where \( \omega \) and \( \omega' \) are symplectic forms on \( M \) and \( M' \) respectively. From (5.2), \( M \) and \( M' \) have the same genus, so that there is a diffeomorphism \( \psi : M \to M' \). Since \( \omega \) and \( \omega' \) are symplectic on \( M \) and \( M' \) and \( H^2_{\text{dr}}(M; \mathbb{R}) \cong \mathbb{R} \cong H^2_{\text{dr}}(M'; \mathbb{R}) \), by Moser’s trick, we can assume that

\[
\psi^* \omega' = T \omega, \quad \text{for some } T > 0.
\]

As both \( p \) and \( p' \) are trivial bundles, it is immediate to find a diffeomorphism

\[
\Psi : \Sigma \to \Sigma \tag{5.3}
\]

lifting \( \psi \) and such that \( \Psi^* p' = p \). Therefore, \( f \) preserves the orientation of \( \Sigma \) and \( \Psi^* \Omega' = T \Omega \).

We now want to describe the connected components of the space of Zoll forms on \( \Sigma \) with fixed cohomology class. In order to do so, we first determine the connected components of the space of oriented \( S^1 \)-bundles \( \mathcal{P}(\Sigma) \) with the help of classical results in low-dimensional topology. As a preliminary observation, we point out that if \( p \in \mathcal{P}(\Sigma) \), then the class \( [p^{-1}(\text{pt})]_{\mathbb{Z}} \in H_1(\Sigma; \mathbb{Z}) \) is primitive, since its intersection number with a global section of \( p \) is equal to 1. We distinguish three cases: \( M = S^2, M = T^2 \), and \( \text{genus}(M) \geq 2 \).

**Case 1:** \( \Sigma \cong S^2 \times S^1 \). We regard \( S^2 \) as the unit sphere in \( \mathbb{R}^3 \). We recall that the mapping class group of orientation-preserving diffeomorphisms of \( S^2 \times S^1 \) is given by

\[
\text{MCG}(S^2 \times S^1) \cong \frac{\mathbb{Z}}{2\mathbb{Z}} [\Psi_1] \oplus \frac{\mathbb{Z}}{2\mathbb{Z}} [\Psi_2].
\]

Here, the generators \( \Psi_1, \Psi_2 : S^2 \times S^1 \to S^2 \times S^1 \) are given by

\[
\Psi_1(q, \phi) := (-q, -\phi), \quad \Psi_2(q, \phi) := (\rho_\phi(q), \phi), \quad \forall (q, \phi) \in S^2 \times S^1,
\]

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where \( \rho_\phi : S^2 \to S^2 \) is the rotation of angle \( 2\pi \phi \) around the north pole. We consider the standard projection \( p_+ : S^1 \times S^2 \to S^2 \) along \( S^1 \) and we define \( p_- := \Psi^*_1 p_+ \).

We claim that \( \Psi(S^2 \times S^1) \) has four connected components containing respectively \( p_+ \), \( \Psi^*_2 p_+ \), \( p_- \) and \( \Psi^*_2 p_- \). To this purpose, we observe that \( [p^*_+1(pt)]_Z = [\Psi^*_2 p^*_+1(pt)]_Z \) and \( [p^*_-1(pt)]_Z = [\Psi^*_2 p^*_-1(pt)]_Z \) are distinct and yield the two primitive homology classes in \( H_1(S^2 \times S^1; Z) \cong Z \). Therefore, we just need to show that \( p_+ \), \( \Psi^*_2 p_+ \) are not in the same connected component. If, by contradiction, there were a path in \( \Psi(S^2 \times S^1) \) from \( p_+ \) to \( \Psi^*_2 p_+ \), then by Lemma 4.7 (ii) there would exist a diffeomorphism \( \Psi_2 : S^2 \times S^1 \to S^2 \times S^1 \) isotopic to \( \Psi_2 \) such that \( p_+ = (\Psi_2')^* p_+ \). This forces \( \Psi_2' \) to be of the form

\[
\Psi_2'(q, \phi) = (\psi(q), \phi'(x, \phi))
\]

where \( \phi \mapsto \phi'(q, \phi) \) is an orientation-preserving diffeomorphism of \( S^1 \), for all \( q \in S^2 \), and \( \psi \) is an orientation-preserving diffeomorphism of \( S^2 \). However, every orientation-preserving diffeomorphisms of \( S^2 \) is isotopic to the identity and the set \( \text{Diff}_+(S^1) \) of orientation-preserving diffeomorphisms of \( S^1 \) is homotopy equivalent to \( S^1 \), so that \( \pi_2(\text{Diff}_+(S^1)) = 0 \). Thus, the map \( \Psi_2' \) would be isotopic to the identity, which is impossible, as it is isotopic to \( \Psi_2 \).

Let \( p \) be an arbitrary element in \( \Psi(S^2 \times S^1) \). We have either \( [p^{-1}(pt)]_Z = [p^*_+1(pt)]_Z \) or \( [p^{-1}(pt)]_Z = [p^*_-1(pt)]_Z \). We have seen in (5.3) that there exists an isomorphism of oriented \( S^1 \)-bundles \( \Psi : S^2 \times S^1 \to S^2 \times S^1 \) preserving the orientation and such that \( \Psi^*_p p_+ = p \).

Therefore, if \( [p^{-1}(pt)]_Z = [p^*_+1(pt)]_Z \), then \( \Psi \) is either isotopic to \( \text{id}_{S^2 \times S^1} \) or to \( \Psi_2 \) and \( p \) is either homotopic to \( p_+ \) or to \( \Psi^*_2 p_+ \); if \( [p^{-1}(pt)]_Z = [p^*_-1(pt)]_Z \), then \( \Psi \) is either isotopic to \( \Psi_1 \) or to \( \Psi_1 \circ \Psi_2 \) and \( p \) is either homotopic to \( p_- \) or to \( \Psi^*_2 p_- \).

**Case 2:** \( \Sigma \cong T^2 \times S^1 = T^3 \). To every orientation-preserving diffeomorphism \( \Psi \) on \( T^3 \), we can associate the induced map in homology

\[
H_1(\Psi) : H_1(T^3; Z) \to H_1(T^3; Z).
\]

As \( H_1(T^3; Z) \cong \mathbb{Z}^3 \), we can identify \( H_1(\Psi) \) with an element of \( \text{SL}(3; \mathbb{Z}) \). Conversely, every \( A \in \text{SL}(3; \mathbb{Z}) \) gives an orientation-preserving diffeomorphism \( \Psi_A : T^3 \to T^3 \) with \( H_1(\Psi_A) = A \).

Moreover, the mapping class group is computed explicitly through the isomorphism

\[
\text{MCG}(T^3) \to \text{SL}(3; \mathbb{Z}), \quad \Psi \mapsto H_1(\Psi). \quad (5.4)
\]

Let \( p_0 : T^3 \cong T^2 \times S^1 \to T^2 \) be the standard projection along \( S^1 \). For every primitive class \( h \in H_1(T^3; Z) \), there exists \( A_h \in \text{SL}(3; Z) \) with \( A_h \cdot h = [p^{-1}(pt)]_Z \), so that the fibres of the oriented \( S^1 \)-bundle

\[
p_h := \Psi^*_A p_0
\]

lie in the homology class \( h \). The map \( A_h \) is not unique. However, if \( A'_h \) is another such map, then \( \Psi^*_A p_0 \) and \( \Psi^*_{A'_h} p_0 \) are equivalent bundles, and up to equivalence, \( p_h \) depends only on \( h \).

On the other hand, if \( p \in \mathcal{P}(T^3) \), we claim that \( p_{[p^{-1}(pt)]_Z} \) and \( p \) are isotopic. Indeed, let \( \Psi : T^3 \to T^3 \) be an orientation-preserving diffeomorphism with \( \Psi^*_p p_0 = p \), as in (5.3). The equality \( H_1(\Psi)([p^{-1}(pt)]_Z) = [p^*_01(pt)]_Z \) implies that \( H_1(\Psi) = A_{[p^{-1}(pt)]_Z} \) so that there holds \( p_{[p^{-1}(pt)]_Z} = \Psi_{H_1(\Psi)} p_0 \). Moreover, since \( \Psi \) and \( \Psi_{H_1(\Psi)} \) are isotopic due to (5.4), \( \Psi_{H_1(\Psi)} p_0 \) is homotopic to \( p \) and the claim is proven.

**Case 3:** \( \Sigma \cong M \times S^1 \) with genus(\( M \)) \( \geq 2 \). Let \( p_+ : M \times S^1 \to M \) be the oriented \( S^1 \)-bundles given by the standard projection and set \( p_- := -p_+ \). The bundles \( p_+ \) and \( p_- \) are
not homotopic since the homotopy classes of their fibres are different. If \( p \) is any element of \( \mathcal{P}(M \times S^1) \), by [Wal67, Satz (5.5)], there exists a diffeomorphism \( \Psi : M \times S^1 \to M \times S^1 \) isotopic to the identity such that either \( \Psi^*p = p_+ \) or \( \Psi^*p = p_- \). Thus, \( p \) is either isotopic to \( p_- \) or to \( p_+ \).

This finishes the description of the connected components of \( \mathcal{P}(\Sigma) \) in the three cases. To determine the connected components of the space of Zoll forms with fixed cohomology class, we consider the map assigning to each Zoll odd-symplectic form its associated bundle

\[
\Psi_Z : Z(\Sigma) \to \mathcal{P}(\Sigma), \quad \Omega \mapsto [\Omega]
\]

If \( p = \Psi_Z(\Omega) \), then \( \text{PD}([\Omega]) \) is a positive multiple of \([p^{-1}(pt)]\) by (1.6). Moreover, the map \( \Psi_Z \) is surjective, since the quotient manifold of any bundle is an orientable surface, and therefore, possesses a positive symplectic form. In particular, if \( p \in \mathcal{P}(\Sigma) \) and \( C \in H^2_{\text{dR}}(\Sigma) \) are such that \( \text{PD}(C) \) is a positive multiple of \([p^{-1}(pt)]\), then there exists \( \Omega \in Z_C(\Sigma) \subset Z(\Sigma) \) such that \( \Psi_Z(\Omega) = p \). This shows that

- if \( M = S^2 \) or \( \mathbb{T}^2 \), then \( Z_C(\Sigma) \) is not empty, for all \( C \neq 0 \);
- if \( M \) has higher genus, then \( Z_C(\Sigma) \) is not empty if and only if \( C \) is a non-zero element in \( p^*(H^2_{\text{dR}}(M)) \).

We fix now some \( C \in H^2_{\text{dR}}(\Sigma) \) for which \( Z_C(\Sigma) \) is non-empty and consider \( \Psi_{Z,C} := \Psi_Z|_{Z_C(\Sigma)} \). As its non-empty fibres are convex, the connected components of \( Z_C(\Sigma) \) correspond through \( \Psi_{Z,C} \) to the connected components of \( \Psi_{Z,C}(Z_C(\Sigma)) \). The latter set is the union of those connected components of \( \mathcal{P}(\Sigma) \), whose elements \( p \) satisfy \( \text{PD}(C) = a[p^{-1}(pt)] \) for some \( a > 0 \). By the discussion above, this shows that:

- If \( M \) is the two-sphere, then the set \( \Psi_{Z,C}(Z_C(\Sigma)) \) has two connected components since \([p_-^{-1}(pt)] = [(\Psi^*_2p_+)^{-1}(pt)] \) and \([p_+^{-1}(pt)] = [(\Psi^*_2p_-)^{-1}(pt)] \);
- If \( M \) has positive genus, then the set \( \Psi_{Z,C} \) is connected.

The proof of the proposition is completed. \( \square \)

# 6 Action of closed two-forms

## 6.1 The action form

**Definition 6.1.** If \( \Omega \) is a closed two-form on \( \Sigma \), an embedded one-periodic curve \( \gamma : S^1 \to \Sigma \) with \( \dot{\gamma} \in \ker \Omega \) is said to be a **closed characteristics** of \( \Omega \). We write \( \mathcal{X}(\Omega) \) for the set of closed characteristics of \( \Omega \), up to orientation-preserving reparametrisations of \( S^1 \).

Fix a free-homotopy class \( \mathfrak{h} \in [S^1, \Sigma] \) and let \( \Lambda_0(\Sigma) \) be the space of one-periodic curves in \( \Sigma \) with class \( \mathfrak{h} \). In what follows, given a pair \((C, \mathfrak{h})\), we shall study a variational principle on \( \Lambda_0(\Sigma) \), which detects, for all \( \Omega \in \mathcal{Z}_C(\Sigma) \), the elements of \( \mathcal{X}(\Omega) \) belonging to \( \Lambda_0(\Sigma) \). To this purpose, we define the action form \( a = a(\Omega) \in \Omega^1(\Lambda_0(\Sigma)) \) by

\[
a_\gamma(\xi) := \int_{S^1} \Omega(\xi(t), \dot{\gamma}(t)) \, dt, \quad \forall \gamma \in \Lambda_0(\Sigma), \forall \xi \in T_\gamma \Lambda_0(\Sigma).
\]
Any $C^1$-path $\{\gamma_r\} \subset \Lambda_h(\Sigma)$ can also be interpreted as a cylinder $\Gamma : [0,1] \times S^1 \to \Sigma$ with $\Gamma(r,t) = \gamma_r(t)$, so that

$$\int_{[0,1] \times S^1} \{\gamma_r\}^* \Omega = \int_{[0,1] \times S^1} \Gamma^* \Omega. \tag{6.1}$$

Furthermore, we have the following classical observation.

**Lemma 6.2.** The action form $a(\Omega)$ is closed. An embedded periodic curve $\gamma \in \Lambda_h(\Sigma)$ is a closed characteristic of $\Omega$ if and only if $a_\gamma(\Omega) = 0$. \(\square\)

If $\Omega$ and $\Omega'$ are forms in $\Xi^2_c(\Sigma)$ with $\Omega' - \Omega = d\alpha$, then there holds

$$a(\Omega') - a(\Omega) = dB_\alpha, \quad B_\alpha(\gamma) := \int_{S^1} \gamma^* \alpha.$$ 

Thus, it makes sense to define $a(C) := [a(\Omega)] \in H^1_{\text{dr}}(\Lambda_h(\Sigma))$. We have a homomorphism

$$H^2_{\text{dr}}(\Sigma) \to H^1_{\text{dr}}(\Lambda_h(\Sigma)), \quad C \mapsto a(C). \tag{6.2}$$

If $a(\Omega)$ admits a primitive functional, then the zeros of $a(\Omega)$ are critical points of the primitive. In the next lemma, we give a criterion ensuring that $a(\Omega)$ is exact on $\Lambda_h(\Sigma)$, in the case that there exists an oriented $S^1$-bundle with fibres $\mathfrak{h}$. Below, we regard classes in $H^2_{\text{dr}}(M)$ as real homomorphisms on $\pi_2(M)$ through the canonical map $\pi_2(M) \to H_2(M;\mathbb{Z})$.

**Lemma 6.3.** Let $C \in H^2_{\text{dr}}(\Sigma)$ and $(p,c) \in \mathcal{X}_C(\Sigma)$ such that $p : \Sigma \to M$ has minus the real Euler class $e \in H^2_{\text{dr}}(M)$ and the oriented $p$-fibres have class $\mathfrak{h} \in [S^1, \Sigma]$. There holds

$$a(C) = 0 \text{ in } H^1_{\text{dr}}(\Lambda_h(\Sigma)) \iff \ker e|_{\pi_2(M)} \subseteq \ker c|_{\pi_2(M)}.$$ 

**Proof.** Let $\omega$ be any element in $\Xi^2_c(M)$ and define $\Omega := p^* \omega$. The cohomology class $a(C)$ is trivial if and only if its integral over any one-periodic curve $\Gamma : S^1 \to \Lambda_h(\Sigma)$ vanishes. Choosing any oriented fibre $\gamma : S^1 \to \Sigma$ such that $[\gamma] = \mathfrak{h}$, we may assume that $\Gamma(0) = \Gamma(1) = \gamma$, up to homotopy, since $a$ is closed and therefore the integral of $a$ depends only on the homotopy class of $\Gamma$. In view of (6.1), $a(C) \in H^1_{\text{dr}}(\Lambda_h(\Sigma))$ is trivial if and only if, for any such $\Gamma$,

$$\int_{[0,1] \times S^1} \Gamma^* \Omega = 0.$$ 

As $p \circ \gamma$ is a constant curve, we think of $\bar{\Gamma} := p \circ \Gamma$ as a map from $S^2$ into $M$ with homotopy class $[\bar{\Gamma}] \in \pi_2(M)$. We have

$$\int_{[0,1] \times S^1} \Gamma^* \Omega = \int_{S^2} \bar{\Gamma}^* \omega = \langle c, [\bar{\Gamma}] \rangle.$$ 

Hence, the lemma follows if we show that the map $\Gamma \mapsto [\bar{\Gamma}]$ is onto $\ker e|_{\pi_2(M)}$. To see that the image $[\bar{\Gamma}]$ is indeed in $\ker e|_{\pi_2(M)}$, we compute

$$\langle e, [\bar{\Gamma}] \rangle = \int_{S^2} \bar{\Gamma}^* \kappa_\eta = \int_{[0,1] \times S^1} \Gamma^* (d\eta) = \int_{S^1} \gamma^* \eta - \int_{S^1} \gamma^* \eta = 0,$$

where $\eta \in \mathcal{K}(p)$ is any connection for $p$. 31
Then, we show that for any \( v : S^2 \to M \) with \( \langle e, [v] \rangle = 0 \), there is \( \Gamma : [0, 1] \times S^1 \to \Sigma \) with \( \Gamma(0, \cdot) = \Gamma(1, \cdot) = \gamma \) such that \( \bar{\Gamma} = p \circ \Gamma = v \). By the naturality of the Euler class, the restriction of \( p \) over \( v \) admits a global section \( \bar{\Upsilon} : S^2 \to \Sigma \). It satisfies \( p \circ \bar{\Upsilon} = v \). Using the quotient map \([0, 1] \times S^1 \to S^2\), we lift \( \bar{\Upsilon} \) to a map \( \bar{\Upsilon}' : [0, 1] \times S^1 \to \Sigma \). Up to modifying \( v \) within its homotopy class, we can assume that \( \bar{\Upsilon}'(0, \cdot) = \bar{\Upsilon}'(1, \cdot) = \gamma(0) \). Finally, we define

\[
\Gamma : [0, 1] \times S^1 \to \Sigma, \quad \Gamma(s, t) = u(t, \bar{\Upsilon}'(s, t)),
\]

where \( u \in \mathcal{U}(p) \). It follows that \( \Gamma(0, \cdot) = \Gamma(1, \cdot) = \gamma \) and \( \bar{\Gamma} = \bar{\Upsilon}' = \bar{\Upsilon} = v \), as required. \( \square \)

**Remark 6.4.** The condition \( \mathfrak{a}(C) = 0 \) in \( H^1_{dR}(\Lambda_h(\Sigma)) \) depends only on \( C \) and \( h \). Therefore, the condition \( \ker e|_{\pi_2(M)} \subseteq \ker e|_{\pi_2(M)} \) is also independent of the chosen pair \( (p, c) \in Z_C(\Sigma) \) with the property that the \( p \)-fibres are in the class \( h \). This last statement can also be seen directly combining Remark 4.4, equation (4.4), and item (i) in Lemma 5.4.

### 6.2 The action functional on a covering space

In view of Lemma 6.3, we cannot expect the existence of a primitive functional of the action form \( \mathfrak{a} \) in general. One standard way to resolve this problem is to find a primitive functional on a suitable covering space of \( \Lambda_h(\Sigma) \). We define a natural covering space in (6.5) below. However, it will have the small disadvantage that in some cases the primitive functional depends on the choice of a one-form \( \alpha \) such that \( \Omega = \Omega_0 + d\alpha \), where \( \Omega_0 \) is a reference two-form in the same cohomology class of \( \Omega \). We overcome this nuisance through the notion of non-degeneracy, which we introduced in Definition 5.3.

As before, let \( C \) be a class in \( H^2_{dR}(\Sigma) \) and \( \mathfrak{P}^0(\Sigma) \) be a connected component of \( \mathfrak{P}(\Sigma) \). We write \( \mathfrak{h} \in [S^1, \Sigma] \) for the class of the oriented fibres of any bundle in \( \mathfrak{P}^0(\Sigma) \). For the rest of this section, we work under the assumption that

\[
Z_C(\Sigma) = \mathfrak{P}^{-1}(\mathfrak{P}_0(\Sigma)) \cap Z_C(\Sigma) \quad \text{is non-empty and non-degenerate},
\]

where \( \mathfrak{P} : Z(\Sigma) \to \mathfrak{P}(\Sigma) \) is the standard projection. We fix a reference pair \( (p_0, c_0) \in Z_C^0(\Sigma) \) with \( p_0 : \Sigma \to M_0 \) and denote by \( e_0 \in H^2_{dR}(M_0) \) minus the real Euler class of \( p_0 \). From (4.3), we have an embedding

\[
j_{p_0} : \Sigma \to \Lambda_h(\Sigma),
\]

where \( j_{p_0}(z) \) is a parametrisation of the oriented \( p_0 \)-fibre passing through \( z \in \Sigma \). By definition, there holds \( j_{p_0}(z) \in \mathcal{X}(\Omega_0) \). We consider the following covering space of \( \Lambda_h(\Sigma) \):

\[
\bar{\Lambda}_h(\Sigma) := \left\{ [\gamma_r] \mid \gamma_0 \in j_{p_0}(\Sigma), \gamma_r \in \Lambda_h(\Sigma), \forall r \in [0, 1] \right\} / \sim,
\]

where \( \{\gamma^0_r\} \sim \{\gamma^1_r\} \), if there is a homotopy \( \{\gamma^s_r\}_{s \in [0,1]} \) such that

\[
\gamma^0_r \in j_{p_0}(\Sigma), \quad \gamma^0_0 = \gamma^1_0 = \gamma^1_1, \quad \forall s \in [0,1].
\]

We denote the elements of \( \bar{\Lambda}_h(\Sigma) \) as \( [\gamma_r] \) so that the covering map is given by

\[
\bar{\Lambda}_h(\Sigma) \to \Lambda_h(\Sigma), \quad [\gamma_r] \mapsto \gamma_1.
\]

We further choose some \( \omega_0 \in \Xi^2_{\mathfrak{a}_0}(M_0) \) and set as reference form

\[
\Omega_0 := p_0^* \omega_0 \in \Xi^2_C(\Sigma).
\]
Let \( \text{Vol} : \Omega^1(\Sigma) \to \mathbb{R} \) and \( \text{Vol}_1 : \Xi^2_C(\Sigma) \to \mathbb{R} \) be the volume functionals associated with \( \Omega_0 \).

As observed in Lemma 5.2 \( \text{Vol}_1 \) depends only on \((p_0, \omega_0)\) but not on \(\omega_0\).

Let \( \alpha \in \Omega^1(\Sigma) \) and recall the notation \( \Omega_\alpha = \Omega_0 + d\alpha \in \Xi^2_C(\Sigma) \). We define the action

\[
\bar{A}_\alpha : \tilde{\Lambda}_b(\Sigma) \to \mathbb{R}, \quad \bar{A}_\alpha([\gamma_\alpha]) := \int_{S^1} \gamma_0^* \alpha + \int_{[0,1] \times S^1} \Gamma^* \Omega_\alpha, \tag{6.7}
\]

where, as before, \( \Gamma : [0,1] \times S^1 \to \Sigma \) is the cylinder associated with \( \{\gamma_\alpha\} \).

The action is well-defined as one sees by integrating \( 0 = d\Omega_\alpha \) over a homotopy satisfying (6.6) and then applying Stokes’ Theorem. Decomposing \( \Omega_\alpha = \Omega_0 + d\alpha \) in the second integrand above and using Stokes’ Theorem, we can rewrite \( \bar{A}_\alpha \) as

\[
\bar{A}_\alpha([\gamma_\alpha]) = \int_{S^1} \gamma_0^* \alpha + \int_{D^2} \Gamma^* \omega_\alpha, \tag{6.8}
\]

where \( \tilde{\Gamma} = p_0 \circ \Gamma : D^2 \to M_0 \). A straightforward computation shows that

\[
d_\gamma \bar{A}_\alpha([\xi]) = a_\gamma(\xi_1), \quad \forall [\gamma_\alpha] \in \tilde{\Lambda}_b(\Sigma), \quad \forall [\xi_\alpha] \in T_{[\gamma_\alpha]} \tilde{\Lambda}_b(\Sigma),
\]

where \( a = a(\Omega_\alpha) \). Hence, Lemma 6.2 can be rephrased as follows.

**Corollary 6.5.** Let \( [\gamma_\alpha] \in \tilde{\Lambda}_b(\Sigma) \) with \( \gamma_1 \in \Lambda_b(\Sigma) \) embedded. Then \( [\gamma_\alpha] \) is a critical point of \( \bar{A}_\alpha \) if and only if \( \gamma_1 \in \mathcal{X}(\Omega_\alpha) \). \( \square \)

If \( \alpha \) and \( \alpha' \) are one-forms such that \( \Omega_\alpha = \Omega = \Omega_{\alpha'} \), then \( \tau := \alpha' - \alpha \) is closed and we have

\[
\bar{A}_{\alpha'} = \bar{A}_\alpha + \langle [\tau], [p^{-1}(\text{pt})] \rangle. \tag{6.9}
\]

**Case 1:** \( e_0 \neq 0 \).

By Lemma 4.5, we have \( [p^{-1}(\text{pt})] = 0 \). The action functional depends only on the two-form \( \Omega \in \Xi^2_C(\Sigma) \) due to (6.9). Therefore, for any \( \alpha \in \Omega^1(\Sigma) \) with \( \Omega = \Omega_\alpha \), we can set

\[
\tilde{A}_\Omega := \bar{A}_\alpha.
\]

In this situation, the non-degeneracy (6.3) of \( \mathcal{Z}^0_C(\Sigma) \) is not needed to associate the action functional with elements in \( \Xi^2_C(\Sigma) \), as opposed to the next case.

**Case 2:** \( e_0 = 0 \).

Here, the action functionals of \( \alpha \) and \( \alpha' \) might be different. Nevertheless, as \( \mathcal{Z}^0_C(\Sigma) \) is non-degenerate, we have \( \langle c_0^\alpha, [M_0] \rangle = \text{ev}(p_0, e_0) \neq 0 \). Thus, by (2.2) and (4.6), we can write

\[
\langle [\tau], [p^{-1}(\text{pt})] \rangle = \frac{\langle [\tau] \cup C_n, [\Sigma] \rangle}{\langle c_0^\alpha, [M_0] \rangle} = \frac{\text{Vol}(\alpha') - \text{Vol}(\alpha)}{\langle c_0^\alpha, [M_0] \rangle},
\]

so that if \( \alpha \) and \( \alpha' \) have the same volume, they also have the same action by (6.9). We set

\[
\tilde{A}_\Omega := \bar{A}_\alpha \tag{6.10}
\]

for a normalised \( \alpha \in \Omega^1(\Sigma) \), i.e. \( \text{Vol}(\alpha) = 0 \), with \( \Omega = \Omega_\alpha \).
Remark 6.6. This remark is parallel to Remark 2.6. For contact forms and Hamiltonian systems, the action functional recovers the following well-known formulae.

- (Contact forms) Let $\Omega_0 = 0$ and $\alpha \in \Omega^1(\Sigma)$ be a (possibly contact) one-form. As $C = 0$, we have that $e_0 \neq 0$ by Lemma 4.5(ii). Thus, this is a special instance of Case 1. In fact, $\alpha(C) = 0$ due to (6.7), and the function $A_\alpha : \Lambda h(\Sigma) \to \mathbb{R}$, given by

$$A_\alpha(\gamma) = \int_{S^1} \gamma^* \alpha,$$

is the unique primitive of $\alpha$ such that $\tilde{A}_{d\alpha}([\gamma_r]) = A_\alpha(\gamma_1)$, for every $[\gamma_r] \in \tilde{\Lambda}_h(\Sigma)$.

- (Hamiltonian systems) Let $p_0$ be trivial, namely $\Sigma = M_0 \times S^1$, so that this is a special instance of Case 2. Assume $\Omega_0 = p_0^* \omega_0$, where $\omega_0$ is a symplectic form on $M_0$, and $\alpha = H dt$, where $t$ is the angular coordinate on $S^1$. If $\gamma_1(t) = (q_1(t), t)$, then we get the classical Hamiltonian action functional

$$\bar{A}_{H dt}([\gamma_r]) := \int_{S^1} H(q_1(t), t) dt + \int_{D^2} \Gamma^* \omega_0$$

on the space of contractible curves with capping disc. Furthermore, the condition $\ker e_0|\pi_2(M_0) \subseteq \ker c_0|\pi_2(M_0)$ in Lemma 6.3 means that $\omega_0$ is symplectically aspherical.

Next, we study the relation between the actions with respect to two different reference weakly Zoll pairs.

Proposition 6.7. Let $(p_0', c_0')$ be another element in $\mathcal{C}_\Omega^0(\Sigma)$. We write $\bar{\Lambda}'_h(\Sigma)$ for the associated covering space of $\Lambda_h(\Sigma)$ and $\tilde{\Lambda}_h(p_0', p_0') \subseteq \tilde{\Lambda}_h(\Sigma)$ for the set of elements $[\delta_r]$ such that $\delta_1 \in \mathcal{P}_0'(\Sigma)$. We pick $\omega_0' \in \Xi^2_{c_0}(M_0')$ and set $\Omega_0' := (p_0')^* \omega_0'$. We choose $\alpha' \in \Omega^1(\Sigma)$ such that

$$\Omega_0' = \Omega_0 + d\alpha', \quad \alpha' \text{ is } \Omega_0\text{-normalised}.$$

For any $\Omega \in \Xi^2_{c_0}(\Sigma)$, we take $\alpha \in \Omega^1(\Sigma)$ such that

$$\Omega = \Omega_0 + d\alpha, \quad \alpha \text{ is } \Omega_0\text{-normalised}.$$

Then, there holds $\Omega = \Omega_0' + d(\alpha - \alpha')$ and the one-form $\alpha - \alpha'$ is $\Omega_0\text{-normalised}$. Moreover, if we denote by $\bar{A}_\Omega$ the action of $\Omega$ with respect to $\Omega_0'$, then

$$\bar{A}_\Omega([\{\delta_r\} \# \{\gamma_r\}]) = \bar{A}_{\Omega_0'}([\delta_r]) + \bar{A}_{\Omega'}([\gamma_r]), \quad (6.11)$$

for every $[\delta_r] \in \tilde{\Lambda}_h(p_0, p_0')$ and $[\gamma_r] \in \tilde{\Lambda}_h(\Sigma)$. Here, the concatenation is made by choosing any representative $\{\gamma_r\}$ of $\{\gamma_r\}$ with $\gamma_0 = \delta_1$.

Proof. The one-form $\alpha - \alpha'$ is $\Omega_0\text{-normalised}$ thanks to Lemma 2.4. Let $\Gamma, \Delta : [0, 1] \times S^1 \to \Sigma$ be the cylinders traced by the paths $\{\gamma_r\}$ and $\{\delta_r\}$, respectively. Let us show equation (6.11) using (6.7) and (6.8):

$$\bar{A}_{\Omega_0'}([\delta_r]) + \bar{A}'_{\Omega'}([\gamma_r]) = \int_{[0,1] \times S^1} \Delta^* \Gamma^* \omega_0 + \int_{[0,1] \times S^1} \gamma_0^* (\alpha - \alpha') + \int_{[0,1] \times S^1} \Gamma^* \omega_0$$

$$= \int_{[0,1] \times S^1} \Delta^* \omega_0 + \int_{[0,1] \times S^1} (\gamma_1)^* \alpha + \int_{[0,1] \times S^1} \Gamma^* \omega_0 + \int_{[0,1] \times S^1} \Gamma^* (d\alpha)$$

$$= \int_{[0,1] \times S^1} (\Delta \# \Gamma)^* \omega_0 + \int_{[0,1] \times S^1} \gamma_0^* \alpha + \int_{S^1} \gamma_1^* \alpha - \int_{S^1} \gamma_0^* \alpha$$

$$= \bar{A}_\Omega([\{\delta_r\} \# \{\gamma_r\}]).$$
Corollary 6.8. The action $\widetilde{A}_\Omega$ depends only on the fixed reference pair $(p_0, c_0) \in \mathcal{Z}_C^0(\Sigma)$, and not on the specific choice of $\omega_0 \in \Xi^2_C(M_0)$.

Proof. Let $\omega'_0$ be another element in $\Xi^2_C(M_0)$. By (6.11) with $(p'_0, c'_0) = (p_0, c_0)$ and $[\delta_r] = [\gamma_0]$, it suffices to show that $\widetilde{A}_{p'_0 \omega'_0}(\gamma_0) = 0$. To this purpose, let $\zeta \in \Omega^1(M_0)$ be such that $\omega'_0 - \omega_0 = d\zeta$. By Lemma 5.2(ii), $p'_0 \zeta$ is $p'_0 \omega_0$-normalised and we can use it to compute $\widetilde{A}_{p'_0 \omega'_0}$:

$$\widetilde{A}_{p'_0 \omega'_0}(\gamma_0) = \int_{S^1} \gamma_0(p'_0 \zeta) = 0.$$

\hfill \Box

6.3 The action is invariant under pull-back and isotopies

In this subsection, we prove two invariance results for the action. For the first one, we consider an additional connected oriented closed manifold $\Sigma^\vee$ of dimension $2n + 1$. We suppose that there are a bundle $p_\vee^0 : \Sigma^\vee \to M_\vee^0$ in $\mathcal{P}^0(\Sigma^\vee)$ and a bundle map $\Pi : \Sigma^\vee \to \Sigma$ with $p_\vee^0 = \Pi^* p_0$. Let $\mathcal{P}^0(\Sigma^\vee)$ be the connected component of $\mathcal{P}(\Sigma^\vee)$ containing $p_\vee^0$. Let $\pi_0 : M_\vee^0 \to M_0$ be the map fitting into the first commutative diagram in (1.1). If we set

$$c_0^\vee := p_0 \pi_0^* c_0 \in H^2_{dR}(M_0), \quad C^\vee := \Pi^* C \in H^2_{dR}(\Sigma^\vee),$$

then $(p_\vee^0, c_0^\vee) \in \mathcal{Z}_{C^\vee}^0(\Sigma^\vee) := (\mathcal{P}^0(\Sigma^\vee))^{-1}(\mathcal{Z}_C^0(\Sigma^\vee)) \cap \mathcal{Z}_{C^\vee}(\Sigma^\vee)$. In particular, $\mathcal{Z}_{C^\vee}^0(\Sigma^\vee)$ is non-empty. We write $\mathcal{P}^0_{C^\vee} : \mathcal{P}^0(\Sigma^\vee) \to \mathcal{P}^0(\Sigma^\vee)$ for the projection and $\text{ev}^\vee : \mathcal{P}^0_{C^\vee}(\Sigma^\vee) \to \mathbb{R}$ for the evaluation. Furthermore, we abbreviate

$$\omega_\vee^0 := \pi_0^* \omega_0 \in \Xi^2_{C^\vee}(M_0^0), \quad \Omega_\vee^0 := (p_\vee^0)^* \omega_\vee^0 = \Pi^* \Omega_0 \in \Xi^2_{C^\vee}(\Sigma^\vee),$$

so that $\Omega_\vee^0$ is associated with the weakly Zoll pair $(p_\vee^0, c_0^\vee)$. We note that

$$\begin{align*}
\pi_0^* (p_0 c_0) &= e_\vee^0, \\
\pi_0^*(A c_0 + c_0) &= A e_\vee^0 + c_0^\vee, \quad \forall A \in \mathbb{R},
\end{align*}$$

(6.12)

where $e_\vee^0$ is minus the real Euler class of $p_\vee^0$. If $h^\vee$ denotes the free-homotopy class of oriented fibres of elements in $\mathcal{P}^0(\Sigma^\vee)$, we write $\widetilde{A}_{h^\vee}(\Sigma^\vee) \to \mathbb{R}$ for the action of $\alpha^\vee \in \Omega^1(\Sigma^\vee)$ with respect to $\Omega_\vee^0$. The bundle map $\Pi$ yields a map between the spaces

$$\Pi_* : \Lambda_{h^\vee}(\Sigma^\vee) \to \Lambda_h(\Sigma), \quad \Pi_* \gamma^\vee := \Pi \circ \gamma^\vee$$

and between their covering spaces

$$\widetilde{\Pi}_* : \Lambda_{h^\vee}(\Sigma^\vee) \to \Lambda_h(\Sigma), \quad \widetilde{\Pi}_*[\gamma^\vee_*] := [\Pi \circ \gamma^\vee_*].$$

Proposition 6.9. The set $\mathcal{Z}_{C^\vee}^0(\Sigma^\vee)$ is non-degenerate if and only if $\deg \pi_0 \neq 0$. In this case, $\widetilde{A}_{h^\vee} : \Lambda_{h^\vee}(\Sigma^\vee) \to \mathbb{R}$ is well-defined for all $\Omega^\vee \in \Xi^2_{C^\vee}(\Sigma^\vee)$ and there holds

$$\widetilde{A}_{h^\vee} \circ \widetilde{\Pi}_* = \widetilde{A}_h \circ \Pi_* \quad \forall \Omega \in \Xi^2_C(\Sigma).$$

Proof. The map $(\mathcal{P}^0_C)^{-1}(p_0) \to (\mathcal{P}^0_{C^\vee})^{-1}(p_\vee^0)$ given by $(p_0, c) \mapsto (p_\vee^0, \pi_0^* c)$ is well-defined. By (1.4) and (6.12), it is also surjective. Since $\mathcal{Z}_C(\Sigma)$ is non-degenerate and there holds

$$\text{ev}^\vee(p_\vee^0, \pi_0^* c) = \deg \pi_0 \cdot \text{ev}(p_0, c), \quad \forall (p_0, c) \in (\mathcal{P}^0_C)^{-1}(p_0),$$

(6.13)

we see that $\mathcal{Z}_{C^\vee}^0(\Sigma^\vee)$ is non-degenerate exactly when $\deg \pi_0$ is non-zero.
Let $\Omega = \Omega_0 + d\alpha \in \mathcal{E}_C^2(\Sigma)$, where $\alpha \in \Omega^1(\Sigma)$ is $\Omega_0$-normalised. Then, $\Pi^* \Omega = \Omega_0^\gamma + d(\Pi^* \alpha)$ and $\Pi^* \alpha$ is $\Omega_0^\gamma$-normalised by Proposition 2.7. For all $[\gamma^\gamma] \in \tilde{\Lambda}_h(\Sigma)$, we have

$$\tilde{A}_{\Pi^* \Omega}([\gamma^\gamma]) = \int_{S^1} (\gamma_0^\gamma)^* \Pi^* \alpha + \int_{[0,1] \times S^1} (\Gamma^\gamma)^* \Pi^* \Omega = \int_{S^1} (\Pi \circ \gamma_0^\gamma)^* \alpha + \int_{[0,1] \times S^1} (\Pi \circ \Gamma^\gamma)^* \Omega$$

$$= \tilde{A}_\Omega(\Pi_* [\gamma^\gamma]).$$

For the second invariance result, we consider a diffeomorphism $\Psi : \Sigma \to \Sigma$ isotopic to $\text{id}_\Sigma$. We define $\Psi^* : \Lambda_h(\Sigma) \to \Lambda_h(\Sigma)$, $\Psi^* \gamma := (\Psi^{-1})_* \gamma = \Psi^{-1} \circ \gamma$.

Let $\{\Psi_r\}$ be an isotopy of diffeomorphisms of $\Sigma$ with $\Psi_0 = \text{id}_\Sigma$ and $\Psi_1 = \Psi$. We denote by $[\Psi_r]$ the homotopy class with fixed end-points of $\{\Psi_r\}$ in the space of isotopies. Given such a homotopy class, we define $[\Psi_r]^* : \Lambda_h(\Sigma) \to \Lambda_h(\Sigma)$, $[\Psi_r]^*[\gamma_r] := [\Psi_r^{-1} \circ \gamma_r]$.

**Proposition 6.10.** For every $\Omega \in \mathcal{E}_C^2(\Sigma)$ and every homotopy class $[\Psi_r]$ as above, there holds

$$\tilde{A}_{\Psi_r^* \Omega} \circ [\Psi_r]^* = \tilde{A}_\Omega.$$

**Proof.** We observe preliminarily that if $X_r$ and $Y_r$ are the time-dependent vector fields generating $\Psi_r$ and $\Psi^{-1}_r$, we have the relation

$$-X_r = d\Psi_r(Y_r) \circ \Psi^{-1}_r.$$

Let $\alpha$ be a normalised one-form such that $\Omega = \Omega_0 + d\alpha$ and let $\{\theta_r\}$ be the path of normalised one-forms introduced in Section 2.2 such that

$$\Psi_r^* \Omega = \Omega_0 + d(\theta_r + \Psi_r^* \alpha), \quad \dot{\theta}_r = \Psi_r^* (\iota_{X_r} \Omega_0), \quad \theta_0 = 0.$$

For $s \in [0, 1]$, we define the truncation $\{\Psi_r^s := \Psi_{rs}\}$ and we write $a_s := a(\Psi_{rs}^s \Omega)$. For every $[\gamma_r] \in \tilde{\Lambda}_h(\Sigma)$, we compute

$$\frac{d}{ds} (\tilde{A}_{\Psi_r^s \Omega}([\Psi_r^s]^*[\gamma_r])) = \left(\frac{d}{ds} \tilde{A}_{\Psi_r^s \Omega}([\Psi_r^s]^*[\gamma_r]) + d_{[\Psi_{rs}^s \gamma_r]} \tilde{A}_{\Psi_r^s \Omega} \cdot \frac{d}{ds} [\Psi_{rs}^s \gamma_r] \right)$$

$$= \int_{S^1} (\Psi_r^s)^* \gamma_s \left(\frac{d}{ds} (\theta_s + \Psi_{rs}^s \alpha)\right) dt + a_{s \gamma_1} \cdot \frac{d}{ds} \Psi_{rs}^s \gamma_1$$

$$= \int_{S^1} \gamma_1 \left(\iota_{X_r} \Omega + d(\alpha(X_r))\right) dt + \int_{S^1} (\Psi_r^s \Omega) \left(\frac{d}{ds} \Psi_{rs}^s \gamma_1, \frac{\partial}{\partial t} \Psi_{rs}^s \gamma_1\right) dt$$

$$= \int_{S^1} \Omega(X_r \circ \gamma_1, \gamma_1) dt + \int_{S^1} \Omega(-X_r \circ \gamma_1, \gamma_1) dt$$

$$= 0,$$

where in the second equality we used (6.8) to compute $\frac{d}{ds} \tilde{A}_{\Psi_r^s \Omega}$. Since $\tilde{A}_{\Psi_0^s \Omega} \circ [\Psi_0]^* = \tilde{A}_\Omega$, the proof is completed. 

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6.4 The action of weakly Zoll pairs

Based on the action functional we have studied, we define the action functional on the space of weakly Zoll pairs:

\[ \mathcal{A} : \mathcal{Z}^0_C(\Sigma) \to \mathbb{R}, \quad \mathcal{A}(p, c) := \tilde{A}_p^*\omega([p_r(z_0)]) \]  

(6.14)

where \( z_0 \) is a point in \( \Sigma \), \( \omega \in \Xi^2_c(M) \), and \( \{ p_r \} \) is a path in \( \mathcal{P}^0(\Sigma) \) starting at the reference bundle \( p_0 \) and ending at \( p_1 = p \). It will turn out that this action is well-defined without the need to pass to a covering space, even when the condition in Lemma 6.3 is not met. This fact is a striking consequence of the non-degeneracy (6.3) of \( \mathcal{Z}^0_C(\Sigma) \). A key role will be played by the following polynomial.

**Definition 6.11.** Let \( Q : \mathbb{R} \to \mathbb{R} \) be the auxiliary polynomial

\[ Q(A) := \text{ev}(p_0, Ae_0 + c_0) = (Ae_0 + c_0)^n, [M_0], \quad \forall A \in \mathbb{R}. \]

The **Zoll polynomial** \( P : \mathbb{R} \to \mathbb{R} \) of the pair \((p_0, c_0)\) is given through

\[ P(0) = 0, \quad \frac{dP}{dA} = Q. \]

**Remark 6.12.** The Zoll polynomial is non-constant by Corollary 5.5 (ii) since we have assumed that \( \mathcal{Z}^0_C(\Sigma) \) is non-degenerate. Furthermore, we have the explicit formula

\[ P(A) = \langle e_0, [M_0] \rangle \frac{A^2}{2} + \langle c_0, [M_0] \rangle A, \quad \text{for } n = 1. \]  

(6.15)

**Remark 6.13.** This remark is parallel to Remark 2.6 and 6.6. The Zoll polynomial has a simple form when any of \( c_0 \) and \( e_0 \) vanishes.

- Let us assume that \( C = 0 \) and take \( c_0 = 0 \), which is relevant to the study of Zoll contact forms. We have

\[ P(A) = \langle e_0^n, [M_0] \rangle \frac{A^{n+1}}{n+1}. \]

- Let us assume that \( e_0 = 0 \), which is relevant to the study of Hamiltonian systems on \( M \) (in which case \( p_0 \) is trivial and \( \omega_0 \) is a symplectic form). We have

\[ P(A) = \langle c_0^n, [M_0] \rangle A. \]

**Theorem 6.14.** The functional \( \mathcal{A} : \mathcal{Z}^0_C(\Sigma) \to \mathbb{R} \) does not depend on any choice involved. Moreover, for any \((p, c) \in \mathcal{Z}^0_C(\Sigma)\), there holds

\[ \mathcal{A}(p, c) = 0, \quad \text{if } e_0 = 0 \quad \text{or } \quad (p, c) = (p_0, c_0), \]  

(6.16)

and

\[ \text{Vol}(p, c) = P(\mathcal{A}(p, c)). \]  

(6.17)

If \( \Psi : \Sigma \to \Sigma \) is isotopic to \( \text{id}_\Sigma \) satisfying \( \Psi^*p = p_0 \) (which exists by Lemma 4.7) and \( \psi : M_0 \to M \) is its quotient map, then

\[ \psi^*c = \mathcal{A}(p, c)e_0 + c_0. \]  

(6.18)
Proof. To ease the notation, we write \((p_1, c_1)\) instead of \((p, c)\) to denote an arbitrary element of \(\mathcal{S}^0_C(\Sigma)\) with \(p_1 : \Sigma \to M_1\) and \(c_1 \in H^2_{dR}(M_1)\). Moreover, \({p_r}\) will indicate any path connecting the reference bundle \(p_0\) with \(p_1\). We divide the proof in five steps.

Step 1. For any path \({p_r}\), the action value \(\tilde{A}_{p_1}^{*}ω_1([p_r(z_0)])\) does not depend on the choice of \(z_0 \in \Sigma\) and \(ω_1 \in \mathcal{E}^2_{C}(M_1)\).

Let \(z_1\) be another point in \(\Sigma\), and let \({z_s}\) be a path between \(z_0\) and \(z_1\) with \(s \in [0, 1]\). We claim that \(\tilde{A}_{p_1}^{*}ω_1([p_r(z_s)])\) does not depend on \(s\). Indeed, since \(\tilde{A}_{p_1}^{*}ω_1 \in \mathcal{X}(p_1^{*}ω_1)\) for all \(s \in [0, 1]\), we see that \(\frac{d}{ds}[p_r(z_s)]\tilde{A}_{p_1}^{*}ω_1 = 0\), by Corollary 6.3 which in turn implies

\[
\frac{d}{ds} \tilde{A}_{p_1}^{*}ω_1([p_r(z_s)]) \bigg|_{s=0} = \frac{d}{ds}[p_r(z_s)]\tilde{A}_{p_1}^{*}ω_1 = 0.
\]

This shows the independence of the action value from \(z_0\).

Next, we take another \(ω_1' \in \mathcal{E}^2_{C}(M_1)\) such that \(ω_1' - ω_1 = dζ\) for some one-form \(ζ\) on \(M_1\). By Lemma 5.2 (ii), \(p_1^{*}ζ\) is normalised with respect to \(p_1^{*}ω_1\). Applying equation (6.11) with \(Ω = p_1^{*}ω_1', Ω_0 = p_1^{*}ω_1\), \({δ_r}\) = \({p_r(z_0)}\) and \({γ_r}\) = \({γ_0}\) a constant path, we find

\[
\tilde{A}_{p_1}^{*}ω_1([p_r(z_0)]) - \tilde{A}_{p_1}^{*}ω_1([p_r(z_0)]) = \tilde{A}_{p_1}^{*}ω_1([γ_0]) = \int_{S^1} γ_0^{*}(p_1^{*}ζ) = 0.
\]

Hence, the action depends on the cohomology class \(c_1\), not on the representative.

Step 2. Let \({p_r}\) be a path and \({Ψ_r}\) an isotopy of diffeomorphisms given by Lemma 4.7 (ii) such that \(Ψ_r^{*}p_r = p_0\). If \(ψ_1 : M_0 \to M_1\) is the quotient map of \(Ψ_1\), then

\[
(i) \quad Ψ_r^{*}c_1 = \tilde{A}_{p_1}^{*}ω_1([p_r(z_0)]) e_0 + c_0, \quad (ii) \quad \tilde{A}_{p_1}^{*}ω_1([p_r(z_0)]) = 0, \quad \text{if} \ e_0 = 0.
\]

Setting \(γ_1 := Ψ_0^{*}p_0 \circ Ψ_r^{-1}(z_0)\) and using Proposition 6.10 we compute

\[
\tilde{A}_{p_1}^{*}ω_1([p_r(z_0)]) = \tilde{A}_{p_1}^{*}ω_1([Ψ_0^{*}p_0 \circ Ψ_r^{-1}(z_0)]) = \tilde{A}_{Ψ_1p_1}^{*}ω_1([Ψ_0^{*}p_0 \circ Ψ_r^{-1}(z_0)]) = \tilde{A}_{Ψ_1p_1}^{*}ω_1([γ_1]) = \int_{S^1} γ_1^{*}(p_1^{*}ζ) = A_1e_0 + c_0.\]

(6.19)

where the last equality follows from \([Ψ_0^{*}p_0 \circ Ψ_r^{-1}(z_0)]) = [γ_1] (see (6.15)). We observe that \(Ψ_0^{*}p_1^{*}ω_1 = p_0^{*}ψ_1^{*}ω_1\) and thus \((p_0, ψ_1^{*}c_1) \in \mathcal{S}^0_C(\Sigma)\). By items (iii) and (iv) in Corollary 5.5 there exists \(A_1 \in \mathbb{R}\) such that

\[
ψ_1^{*}c_1 = A_1e_0 + c_0.\]

(6.20)

If \(e_0 \neq 0\), \(A_1\) is uniquely defined by this property. If \(e_0 = 0\), we simply impose \(A_1 := 0\). Equation (6.20) implies that there are \(η \in \mathcal{K}(p_0)\) and \(ζ \in Ω^1(M_0)\) such that

\[
Ψ_1^{*}p_1^{*}ω_1 = p_0^{*}ω_0 + d(A_1 η + p_0^{*}ζ), \quad p_0^{*}κ = dη
\]

for some \(κ \in \mathcal{E}^2_{C}(M_0)\). Since \(A_1 = 0\), if \(e_0 = 0\), the one-form \(A_1 η + p_0^{*}ζ\) is \(p_0^{*}ω_0\)-normalised by Lemma 5.2 (ii) and we conclude with (6.19) that

\[
\tilde{A}_{p_1}^{*}ω_1([p_r(z_0)]) = \tilde{A}_{Ψ_1p_1}^{*}ω_1([γ_1]) = \int_{S^1} γ_1^{*}(A_1 η + p_0^{*}ζ) = A_1.
\]

This proves both items in Step 2.
Step 3. For any path \{p_r\}, there holds \( \mathfrak{Vol}(p_1, c_1) = P(\tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\})) \).

From item (iii) in Lemma 5.4 with \( A = 0 \), we have

\[ \mathfrak{Vol}(p_1, c_1) = \mathfrak{Vol}(p_0, \psi_1^* c_1) = \text{Vol}(A_1 \eta + \psi_0^* \zeta), \]

where \( A_1 = \tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\}) \), \( \eta \) and \( \zeta \) are from the previous step. Using the definition of the volume (2.1) and Fubini’s Theorem, we compute

\[
\text{Vol}(A_1 \eta + \psi_0^* \zeta) = \int_0^1 \left( \int_{\Sigma} (A_1 \eta + \psi_0^* \zeta) \wedge \psi_0^* (rA_1 \kappa + \omega_0)^n \right) dr
\]

\[
= \int_0^1 \left( \int_{M_0} ((p_0)^*(A_1 \eta + \psi_0^* \zeta)) \wedge (rA_1 \kappa + \omega_0)^n \right) dr
\]

\[
= \int_0^1 A_1( (rA_1 e_0 + c_0)^n, [M_0]) dr
\]

\[
= \int_0^1 A_1 Q(rA_1) dr
\]

\[
= P(A_1) - P(0)
\]

\[
= P(A_1).
\]

Step 4. If \( (p_1, c_1) = (p_0, c_0) \), then, for any path \{p_r\}, we have \( \tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\}) = 0 \).

Let \( \psi_1 \) and \( A_1 = \tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\}) \) as in Step 2. Since \( \psi_1^* e_0 = e_0 \), applying \( \psi_1 \) iteratively to item (i) in Step 2, we obtain

\[
(\psi_1^m)^* e_0 = mA_1 e_0 + c_0, \quad \forall m \in \mathbb{Z}.
\]

From item (ii) in Lemma 5.4 with \( A = 0 \), we deduce

\[
Q(0) = \text{ev}(p_0, c_0) = \text{ev}(p_0, (\psi_1^m)^* e_0) = \text{ev}(p_0, mA_1 e_0 + c_0) = Q(mA_1), \quad \forall m \in \mathbb{Z}.
\]

We assume by contradiction that \( A_1 \neq 0 \). Since \( Q \) is a polynomial, the above identity implies that \( Q(A) \equiv Q_0 \), where \( Q_0 \in \mathbb{R} \), so that \( P(A) = Q_0 A \) for all \( A \in \mathbb{R} \). Since \( \mathfrak{Z}_C(\Sigma) \) is non-degenerate, the coefficient \( Q_0 \) is non-zero. This contradicts Step 3:

\[
0 = \mathfrak{Vol}(p_0, c_0) = P(A_1) = Q_0 A_1.
\]

Step 5. End of the proof.

To show that the functional \( \mathcal{A} : \mathfrak{Z}_C(\Sigma) \rightarrow \mathbb{R} \) is well-defined, it remains to see that the action value \( \tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\}) \) does not depend on the choice of path \{p_r\}. Indeed, if \( \{p_r'\} \) is another path with \( p_0' = p_0 \) and \( p_1' = p_1 \), the concatenation \( \{p_r'\} \# \{p_{1-r} \} \) forms a loop based at \( p_0 \). Applying Step 4 and the definition of the action (6.8), we conclude

\[
\tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\}) = \tilde{A}_{p_1 \omega_1}(\{p_r(z_0)\}) = \tilde{A}_{p_1 \omega_0}(\{p_r' \# p_{1-r}(z_0)\}) = 0.
\]

Finally, the identities claimed in the statement of the theorem follow directly from what we have proven. Equation (6.16) is a consequence of item (ii) in Step 2 and of Step 4. Equation (6.17) follows from Step 3. Item (i) in Step 2 implies (6.18).
Remark 6.15. If $\varepsilon = 0$, we could also have considered the space of weakly Zoll one-forms $\mathfrak{z}^0_{C}(\Sigma)$, whose elements are pairs $(p, \alpha)$ with $(p, [\omega]) \in \mathfrak{z}^0_{C}(\Sigma)$, where $p^*\omega_\alpha = \Omega + d\alpha$, and $\alpha$ not necessarily normalised. We have an action functional

$$\mathcal{A} : \mathfrak{z}^0_{C}(\Sigma) \to \mathbb{R}, \quad \mathcal{A}(p, \alpha) := \mathcal{A}_\alpha([j_p(z_0)]).$$

Since $\mathfrak{z}^0_{C}(\Sigma)$ is non-degenerate, from Remark 6.13 and Theorem 6.14, we deduce

$$\text{Vol}(\alpha) = \langle c^\alpha_0, [M_0]\rangle \cdot \mathcal{A}(p, \alpha), \quad \forall (p, \alpha) \in \mathfrak{z}^0_{C}(\Sigma).$$

Remark 6.16. Let $\Omega \in Z_{C}(\Sigma)$ be a Zoll odd-symplectic form such that the associated bundle $p_\Omega$ belongs to $\mathfrak{z}^0(\Sigma)$. If we define

$$\mathcal{A}(\Omega) := \mathcal{A}(p_\Omega, [\omega_\Omega]), \quad \Omega = p^*_\Omega \omega_\Omega,$$

then Theorem 6.14 translates into

$$\text{Vol}(\Omega) = P(\mathcal{A}(\Omega)), \quad \forall \Omega \in Z_{C}(\Sigma).$$

As we did for the volume and the action, we study how the Zoll polynomial behaves under pull-back and change of pair $(p_0, c_0)$ within $\mathfrak{z}^0_{C}(\Sigma)$.

Proposition 6.17. Let $p'_0 : \Sigma^\vee \to M'_0$ be an oriented $S^1$-bundle with maps $\Pi : \Sigma^\vee \to \Sigma$ and $\pi_0 : M'_0 \to M_0$ as described in (6.11). If $P'$ is the Zoll polynomial of the weakly Zoll pair $(\Pi^* p_0, \pi_0^* c_0) = (p'_0, c'_0)$, there holds

$$P' = (\text{deg} \pi_0) \cdot P.$$

Proof. From equations (6.12) and (6.13), we see that $Q' = (\text{deg} \pi_0) \cdot Q$, where $Q'$ is the derivative of $P'$. Since $P(0) = 0 = P'(0)$, the statement follows.

Proposition 6.18. Let $(p'_0, c'_0)$ be a pair in $\mathfrak{z}^0_{C}(\Sigma)$ with $p'_0 : \Sigma \to M'_0$. There holds

$$\langle (c'_0)^n, [M'_0]\rangle = \text{ev}(p'_0, c'_0) = \frac{dP}{dA}(A(p'_0, c'_0)).$$

If $P'$ is the Zoll polynomial associated with $(p'_0, c'_0)$, then

$$P'(A) = P(A + \mathcal{A}(p'_0, c'_0)) - P(A(p'_0, c'_0)), \quad \forall A \in \mathbb{R}.$$

Proof. If $Q'$ denotes the derivative of $P'$, by Lemma 5.4(ii) and (6.18), we have

$$Q'(A) = Q(A + \mathcal{A}(p'_0, c'_0)).$$

Setting $A = 0$, we get the first part of the statement. For the second part, we integrate the above identity and use the normalization $P'(0) = 0$:

$$P'(A) = \int_0^A Q'(A')dA' = \int_0^A Q(A' + \mathcal{A}(p'_0, c'_0))dA' = P(A + \mathcal{A}(p'_0, c'_0)) - P(A(p'_0, c'_0)).$$

Combining the result above with the transformation rules for the action and the volume under change of reference weakly Zoll pair, we conclude the following invariance property.
Corollary 6.19. Let $P'$, $\tilde{A}$, $\Psi\Omega'$ be the Zoll polynomial, the action and the volume associated with another reference pair $(p'_0, c'_0) \in \mathcal{Z}_C^0(\Sigma)$. For every $\Omega \in \operatorname{Zol}(\Sigma)$, there holds
\[
P' \circ \tilde{A}'_\Omega - \Psi\Omega'(\Omega) = P \circ \tilde{A}_\Omega - \Psi\Omega(\Omega) \quad \text{on} \quad \tilde{A}_h(\Sigma).
\]

Proof. We preliminarily observe that
\[
\tilde{A}_\Omega = A(p'_0, c'_0) + \tilde{A}'_\Omega
\]
thanks to Proposition 6.7. We also recall that Lemma 2.4 implies that
\[
\Psi\Omega(\Omega) = \Psi\Omega(p'_0, c'_0) + \Psi\Omega'(\Omega).
\]

Using these identities together with Proposition 6.18 and Theorem 5.14, we readily compute
\[
P(\tilde{A}_\Omega) - \Psi\Omega(\Omega) = P(\tilde{A}_\Omega) - \Psi\Omega(p'_0, c'_0) - \Psi\Omega'(\Omega)
\]
\[
= P(\tilde{A}_\Omega + A(p'_0, c'_0)) - \Psi\Omega(p'_0, c'_0) - \Psi\Omega'(\Omega)
\]
\[
= P'(\tilde{A}'_\Omega) + P(A(p'_0, c'_0)) - \Psi\Omega(p'_0, c'_0) - \Psi\Omega'(\Omega)
\]
\[
= P'(\tilde{A}'_\Omega) - \Psi\Omega'(\Omega). \quad \square
\]

7 A conjectural local systolic-diastolic inequality

The following is our setting to study a systolic-diastolic inequality for odd-symplectic forms. We assume that there exists a Zoll odd-symplectic form $\Omega_0 \in \mathcal{Z}_C(\Sigma)$, for some $C \in H^2_{dR}(\Sigma)$. We denote the bundle associated with $\Omega_0$ by
\[
p_1 := p_{\Omega_0} : \Sigma \to M_1 := M_{\Omega_0}
\]
and let $\omega_0 \in \mathcal{Z}_C^2(M_1)$ be the symplectic form such that $\Omega_0 = p_1^*\omega_0$. The space $\mathcal{Z}_C^0(\Sigma)$ is the connected component of $p_1$ in $\mathcal{Z}^2_{\Sigma}(\Sigma)$ and $\mathcal{Z}_C^0(\Sigma)$ is the corresponding space of weakly Zoll pairs, which is non-empty and non-degenerate thanks to Remark 5.7. Namely, condition (6.3) is automatically satisfied and we can use all the results contained in Section 6.

We fix a reference weakly Zoll pair $(p_0, c_0) \in \mathcal{Z}_C^0(\Sigma)$, where
\[
p_0 : \Sigma \to M_0
\]
and let $\omega_0$ be some form in $\mathcal{Z}_{c_0}^2(M_0)$. Let $h \in [S^1, \Sigma]$ denote the free-homotopy class of the $p_0$-fibres. We define the volume and the action functionals
\[
\Psi\Omega : \mathcal{Z}_{c_0}^2(\Sigma) \to \mathbb{R}, \quad A : \mathcal{Z}_C^0(\Sigma) \to \mathbb{R}, \quad \tilde{A}_\Omega : \tilde{A}_h(\Sigma) \to \mathbb{R}
\]
with respect to $(p_0, c_0)$. Due to Proposition 6.18 we have
\[
\frac{dP}{dA}(A_*) = ([\omega_0]^n, [M_1]) > 0, \quad A_* := A(p_1, \omega_0).
\]

Definition 7.1. We say that a finite cover $\{B_i\}$ of $M_1$ is admissible, if all the sets $B_i$ and their intersections $B_i \cap B_j$ are homeomorphic to the open Euclidean ball, if non-empty. We write $\{\Sigma_i := p_1^{-1}(B_i)\}$ for the pull-back cover of $\Sigma$, so that $\Sigma_i \cap \Sigma_j$ retracts to a $p_1$-fibre, if non-empty. We readily see that admissible finite covers exist and we fix one of them throughout the present section.
We also consider another oriented \( S^1 \)-bundle \( p_0^\gamma : \Sigma^\gamma \to M_0^\gamma \) with bundle map \( \Pi : \Sigma^\gamma \to \Sigma \) such that \( p_0^\gamma = \Pi^\gamma p_0 \). We further assume that \( \Pi \) is a **finite covering map**. This is equivalent to asking that the corresponding quotient map \( \pi_0 : M_0^\gamma \to M_0 \), which fits into the first commutative diagram in (4.11), is a finite covering map. Thus, we endow \( M_0^\gamma \) and \( \Sigma^\gamma \) with the orientations pulled back by \( \pi_0 \) and \( \Pi \), respectively.

The pull-back form \( \Omega^\gamma_* := \Pi^\gamma \Omega_* \) is Zoll odd-symplectic on \( \Sigma^\gamma \) and defines the oriented \( S^1 \)-bundle \( p_1^\gamma := \Pi^\gamma p_1 : \Sigma^\gamma \to M_1^\gamma \). Since \( p_0 \) and \( p_1 \) are connected by a path of oriented \( S^1 \)-bundles, the quotient map \( \pi_1 : M_1^\gamma \to M_1 \) is also an orientation-preserving finite covering map, which fits into the first commutative diagram in (4.11). Moreover, \( \pi_1 \) evenly covers all the sets in \( \{ B_i \} \), as they are contractible. Therefore, the family \( \{ B_i \} \), whose elements are the connected components of the pre-images \( \pi_1^{-1}(B_i) \), yields a finite admissible cover of \( M_1^\gamma \).

### 7.1 A local primitive for the action form

We have seen that the action form does not admit a global primitive on \( \Lambda_0(\Sigma) \) in general. However, we can always find a local primitive on the space of periodic curves that are close to \( j_{p_1}(\Sigma) \). Let \( \Lambda(p_1) \) be the subset of \( \Lambda_0(\Sigma) \), whose elements \( \gamma \) are contained in some \( \Sigma_i \) and are homotopic within \( \Sigma_i \) to some \( p_1 \)-fibre. This means that \( \gamma \) admits a **short homotopy** \( \{ \gamma_r^\text{short} \} \). Namely, such a homotopy has the properties

\[
\gamma_r^\text{short} \subset \Sigma_i, \quad \forall r \in [0, 1], \quad \gamma_0^\text{short} = j_{p_1}(z_i), \quad \gamma_1^\text{short} = \gamma,
\]

for some \( z_i \in \Sigma_i \). We choose any path \( \{ p_r \} \) from \( p_0 \) and \( p_1 \) in \( \mathcal{P}^0(\Sigma) \), and consider the map

\[
\Lambda(p_1) \to \tilde{\Lambda}_0(\Sigma), \quad \gamma \mapsto \left[ \{ j_{p_r}(z_i) \} \# \{ \gamma_r^\text{short} \} \right]
\]

Since the intersection of two elements of the finite admissible cover is contractible, this map does not depend on the particular short homotopy nor on the set \( \Sigma_i \). Composing this map with the functional \( \tilde{A}_\Omega \) on \( \tilde{\Lambda}_0(\Sigma) \), we define the action functional on \( \Lambda(p_1) \):

\[
A_\Omega : \Lambda(p_1) \to \mathbb{R}, \quad A_\Omega(\gamma) := \tilde{A}_\Omega\left( \left[ \{ j_{p_r}(z_i) \} \# \{ \gamma_r^\text{short} \} \right] \right). \tag{7.2}
\]

We note that \( dA_\Omega = \alpha(\Omega) \) on \( \Lambda(p_1) \).

**Lemma 7.2.** The action value \( A_\Omega(\gamma) \) does not depend on the choice of path \( \{ p_r \} \).

*Proof.* If \( \{ p'_r \} \) is another path in \( \mathcal{P}^0(\Sigma) \) with \( p'_0 = p_0 \) and \( p'_1 = p_1 \), then, by (6.8) and (6.16),

\[
\tilde{A}_\Omega\left( \left[ \{ j_{p_r}(z_i) \} \# \{ \gamma_r^\text{short} \} \right] \right) - \tilde{A}_\Omega\left( \left[ \{ j_{p'_r}(z_i) \} \# \{ \gamma_r^\text{short} \} \right] \right) = A(p_0, c_0) = 0. \quad \square
\]

We write \( \mathcal{X}(\Omega; p_1) := \mathcal{X}(\Omega) \cap \Lambda(p_1) \) for the set of closed characteristics of \( \Omega \) in \( \Lambda(p_1) \). The **systole** over the set \( \Lambda(p_1) \) is the functional

\[
A_{\text{min}} : \Xi^2_c(\Sigma) \to \mathbb{R} \cup \{-\infty, +\infty\}, \quad A_{\text{min}}(\Omega) := \inf_{\gamma \in \mathcal{X}(\Omega; p_1)} A_\Omega(\gamma), \tag{7.3}
\]

while the **diastole** over the set \( \Lambda(p_1) \) is the functional

\[
A_{\text{max}} : \Xi^2_c(\Sigma) \to \mathbb{R} \cup \{-\infty, +\infty\}, \quad A_{\text{max}}(\Omega) := \sup_{\gamma \in \mathcal{X}(\Omega; p_1)} A_\Omega(\gamma). \tag{7.4}
\]
Proposition 7.3. Let $\Pi : \Sigma^\gamma \to \Sigma$ be a bundle map with $p_0^\gamma = \Pi^* p_0$, which is also a finite covering map. Let us consider the oriented $S^1$-bundle $p_1^\gamma = \Pi^* p_1 : \Sigma^\gamma \to M_1^\gamma$. Then, $\Pi_* : (\gamma^\gamma \in \Lambda(p_1^\gamma)) \mapsto (\Pi \circ \gamma^\gamma \in \Lambda(1))$ is a (surjective) finite covering map and there holds

$$A_{\Pi^* \Omega} = A_\Omega \circ \Pi_* \quad \forall \Omega \in \mathcal{E}_C^2(\Sigma),$$

where $A_{\Pi^* \Omega}$ is the action functional on $\Lambda(p_1^\gamma)$ with reference pair $(p_0^\gamma, c_0^\gamma)$. As a consequence, the restriction map $\Pi_* : \mathcal{X}(\Pi^\gamma \Omega; p_1^\gamma) \to \mathcal{X}(\Omega; p_1)$ is also a finite cover and we have

$$A_{\min(\Pi^\gamma \Omega)} = A_{\min(\Omega)}, \quad A_{\max(\Pi^\gamma \Omega)} = A_{\max(\Omega)}.$$

Proof. The statement follows from Proposition 6.9 since $\Pi^*$ lifts a path $\{p_r\}$ in $\mathcal{P}^0(\Sigma)$ to a path $\{p_r^\gamma\}$ in $\mathcal{P}^0(\Sigma^\gamma)$ and $\Pi$ maps short homotopies for $p_1^\gamma$ to short homotopies for $p_1$.

We recall some results from [Gin87], ensuring that the systole and the diastole of $\Omega$ are finite, if $\Omega$ is an odd-symplectic form sufficiently close to $\Omega_*$ in the same cohomology class. As usual, we measure distances and norms by fixing Riemannian metrics on the relevant spaces.

Proposition 7.4. There exists a $C^1$-neighbourhood $\mathcal{V}$ of $\Omega_*$ in $\mathcal{S}_C(\Sigma)$ such that, for all $\Omega \in \mathcal{V}$, the following properties hold:

(i) There is a constant $C > 0$ depending on $\mathcal{V}$ but not on $\Omega$ such that

$$\sup_{t \in \mathbb{R}} \left( \text{dist}(q_\gamma(0), q_\gamma(t)) \right) \leq C \|\Omega - \Omega_*\|_{C^0}, \quad q_\gamma := p_1 \circ \gamma, \quad \forall \gamma \in \mathcal{X}(\Omega; p_1).$$

(ii) The set $\mathcal{X}(\Omega; p_1)$ is compact and non-empty and the restrictions $A_{\min} : \mathcal{V} \to \mathbb{R}$, $A_{\max} : \mathcal{V} \to \mathbb{R}$ are Lipschitz-continuous real functions.

(iii) If $\Omega \in \mathcal{V}$ is a Zoll form, then all its closed characteristics (up to orientation) lie in $\Lambda(p_1)$, i.e. $\mathcal{X}(\Omega) = \mathcal{X}(\Omega; p_1)$, and, recalling the notation $A(\Omega) = A(p_1, [\omega_\Omega])$,

$$A_{\min}(\Omega) = A(\Omega) = A_{\max}(\Omega).$$

Proof. Let $\eta \in \mathcal{K}(p_1)$, and let $V$ be the vertical vector field with $\eta(V) = 1$. If $\Omega$ is odd-symplectic, the equation

$$\iota_{X_\Omega}(\eta \wedge \Omega^n) = \Omega^n$$

determines a nowhere vanishing section $X_\Omega \in \text{ker} \Omega$. The periodic orbits of the flow $\Phi^{X_\Omega}$ yield elements in $\mathcal{X}(\Omega)$. Since $V = X_\Omega$, we have

$$\|X_\Omega - V\|_{C^1} \leq C_1 \|\Omega - \Omega_*\|_{C^1}$$

for some constant $C_1 > 0$. Using this estimate, one can argue as in [BK19a, Proposition 3.4] to show that there exists a $C^1$-neighbourhood $\mathcal{V} \subset \mathcal{S}_C(\Sigma)$ of $\Omega_*$ such that, for all $\Omega \in \mathcal{V}$ item (i) holds and $T(\gamma) \leq 2$, where $T(\gamma)$ is the period of $\gamma \in \mathcal{X}(\Omega; p_1)$.

One can follow [Gin87, Section III] or [ ´APB14, Section 3.2] to show that there exists a Lipschitz-continuous map $\mathcal{Y} : \mathcal{V} \times \Sigma \to \Lambda(p_1)$ with the following properties. If one defines, for every $\Omega \in \mathcal{V}$, the function $S(\Omega) : \Sigma \to \mathbb{R}$ by $S(\Omega)(z) := A_\Omega(\mathcal{Y}(\Omega, z))$ for all $z \in \Sigma$, then one gets a Lipschitz function $S : \mathcal{V} \to C^1(\Sigma)$ such that

$$\text{Crit } S(\Omega) \to \mathcal{X}(\Omega; p_1), \quad z \mapsto \mathcal{Y}(\Omega, z)$$
is a one-to-one correspondence. In particular,
\[ \min S(\Omega) = A_{\min}(\Omega), \quad \max S(\Omega) = A_{\max}(\Omega). \]
Since the functions \( \max, \min : C^1(\Sigma) \to \mathbb{R} \) are Lipschitz-continuous, this shows (ii).

To prove (iii), let \( \rho_{\text{Leb}} > 0 \) be a Lebesgue number for the open cover \( \{B_i\} \) of \( M_1 \) and let \( C \) be the constant given in (i). Up to shrinking \( \mathcal{V} \), we can assume that
\[ C\|\Omega - \Omega_*\|_{C^1} \leq \rho_{\text{Leb}}, \quad \forall \Omega \in \mathcal{V}. \]
Let now \( \Omega \in \mathcal{V} \) be a Zoll form with associated bundle \( p_\Omega \), and consider the set
\[ \Sigma(\Omega; p_1) := \{ z \in \Sigma \mid \{ t \mapsto \Phi_t^{X_\Omega}(z) \} \in \Lambda(p_1) \}, \]
which is non-empty since we proved that \( \mathcal{X}(\Omega; p_1) \) is non-empty. This set is open, as the sets \( B_i \) are open. Furthermore, from item (i) and the inequality \( C\|\Omega - \Omega_*\|_{C^1} \leq \rho_{\text{Leb}} \), it is also closed. Since \( \Sigma \) is connected, we deduce \( \Sigma(\Omega; p_1) = \Sigma \) as desired.

To prove the equality between the actions in (iii), we observe that, since \( j_{q_r}(z_i) \) are contained in \( \Sigma_i \), and, hence, build a short path. Therefore,
\[ \mathcal{A}_\Omega(j_{p\# q_r}(z_i)) = \tilde{A}_\Omega([j_{p\# q_r}(z_i)]) = \tilde{A}_\Omega([j_{p\# q_r}(z_i)]) = \mathcal{A}(\Omega), \]
where \( \{ (p\# q_r) \} = \{ p_r \} \neq \{ q_r \} \). As far as the existence of \( \{ q_r \} \) is concerned, the first statement in Lemma 4.7 yields a diffeomorphism \( \Psi : \Sigma \to \Sigma \), which is \( C^1 \)-close to \( \text{id}_\Sigma \) and satisfies \( \Psi p_\Omega = p_1 \). As a consequence, there exists an isotopy \( \{ \Psi_r \} \) of diffeomorphisms \( C^1 \)-close to \( \text{id}_\Sigma \) with \( \Psi_1 = \Psi \). The path \( q_r := (\Psi_{-1})^* p_0 \) has the desired properties.

It seems appropriate at this point to recall Conjecture 1 formulated in the Introduction.

**Conjecture 2.** If \( \Omega_* \in Z_C(\Sigma) \) is a Zoll odd-symplectic form, there is a \( C^{k-1} \)-neighbourhood \( \mathcal{U} \) of \( \Omega_* \) in \( S_C(\Sigma) \), for some \( k \geq 2 \), such that
\[ P(A_{\min}(\Omega)) \leq \mathfrak{Vol}(\Omega) \leq P(A_{\max}(\Omega)), \quad \forall \Omega \in \mathcal{U}. \]
The equality holds in any of the two inequalities above if and only if \( \Omega \) is Zoll.

The order of the three terms in the inequality stems from the fact that \( P \) is monotone increasing in a neighbourhood of \( A_* \) by (7.1). Theorem 6.14 and Proposition 7.4(iii) imply that if \( \Omega \) is Zoll, then both equalities hold.

**Corollary 5.** If \( \mathcal{V} \subset S_C(\Sigma) \) is the \( C^1 \)-neighbourhood of a Zoll odd-symplectic form \( \Omega_* \) given by Theorem 7.3, then there holds
\[ P(A_{\min}(\Omega)) = \mathfrak{Vol}(\Omega) = P(A_{\max}(\Omega)), \quad \forall \Omega \in \mathcal{V} \cap Z_C(\Sigma). \]

**Remark 6.** If the real Euler class vanishes, then the inequality in Conjecture 2 becomes
\[ A_{\min}(\Omega) \leq 0 \leq A_{\max}(\Omega), \quad \forall \Omega \in \mathcal{U}. \]
Indeed, in this case \( \mathfrak{Vol}(\Omega) \equiv 0 \), \( P(A) = ([\omega_*]^n, M_1] \) and \( ([\omega_*]^n, M_1] > 0 \) by (7.1).
Remark 7.7. The conjecture is independent of the chosen reference pair \((p_0, c_0) \in \mathcal{Z}_C^0(\Sigma)\). Indeed, if \((p_0', c_0') \in \mathcal{Z}_C^0(\Sigma)\) is another pair with corresponding objects \(\mathfrak{vol}', P'\) and \(A'_\Omega\), then

\[
P' \circ A'_{\Omega} - \mathfrak{vol}'(\Omega) = P \circ A_{\Omega} - \mathfrak{vol}(\Omega)
\]

by Corollary 6.19.

In the next subsection, we will see that \(P\) can be written as the integral over \(M_1\) of a function \(P : M_1 \times \mathbb{R} \to \mathbb{R}\) so that the monotonicity of \(P\) corresponds to the monotonicity of \(P\) in the second variable for an interval \(I\) containing \(A_\ast\). This will enable us to establish the conjecture in some simple cases. Before doing that, we end this subsection by observing that Conjecture 1 is also consistent with the pull-back operation.

Proposition 7.8. Let \(\Pi : \Sigma^\vee \to \Sigma\) be a bundle map with \(p_0^\vee = \Pi^*p_0\), which is also a finite covering map. If Conjecture 1 holds for \(\Omega^\vee = \Pi^*\Omega_\ast \in \mathcal{S}_C^\vee(\Sigma^\vee)\), then it holds for \(\Omega_\ast \in \mathcal{S}_C(\Sigma)\).

Proof. Let us suppose that the conjecture is true for a neighbourhood \(U^\vee \subset V^\vee\) around \(\Omega^\vee\), where \(V^\vee\) is given by Proposition 7.3. Thus, for all \(\Omega^\vee \in U^\vee\), there holds

\[
P^\vee(A_{\min}(\Omega^\vee)) \leq \mathfrak{vol}(\Omega^\vee) \leq P^\vee(A_{\max}(\Omega^\vee))
\]

(7.5)

and any of the two equalities holds if and only if \(\Omega^\vee\) is Zoll. We claim that the local systolic inequality holds in the neighbourhood \(U := (\Pi^*)^{-1}(U^\vee)\) of \(\Omega_\ast\). Indeed, let us take an arbitrary \(\Omega\) in \(U\). This means that \(\Pi^*\Omega\) belongs to \(U^\vee\). Applying Proposition 7.3, Proposition 6.17 and Proposition 2.7, it follows that

\[
P(A_{\min}(\Omega)) = P(A_{\min}(\Pi^*\Omega)) = \frac{1}{\deg \pi_0} P^\vee(A_{\min}(\Pi^*\Omega)) \leq \frac{1}{\deg \Pi} \mathfrak{vol}(\Pi^*\Omega) = \mathfrak{vol}(\Omega),
\]

where we used \(\deg \Pi = \deg \pi_0 > 0\) and (7.3). An analogous inequality holds for \(A_{\max}(\Omega)\).

Let us assume now, for instance, that \(P(A_{\min}(\Omega)) = \mathfrak{vol}(\Omega)\). From the computation above, this implies \(P^\vee(A_{\min}(\Pi^*\Omega)) = \mathfrak{vol}(\Pi^*\Omega)\). Since the conjecture is true for \(U^\vee\), we see that \(\Pi^*\Omega\) is Zoll and we call \(p^\vee\) the associated bundle. All the fibres of \(p^\vee\) lie in some \(\Sigma^\vee\) by Proposition 7.4. Since \(\Pi\) is injective on such sets, it follows that there exists a bundle \(p : \Sigma \to M\) such that \(p^\vee = \Pi^*p\), which implies that \(\Omega\) is a Zoll form with bundle \(p\).

\(\square\)

7.2 A second look at H-forms

For the rest of this section, we assume that \(p_0 = p_1 = p_\Omega\), and take \(\{p_r\}\) as the constant path. This does not cause any loss in generality thanks to Remark 7.7. We fix a free \(S^1\)-action \(u \in \mathfrak{u}(p_0)\) with generating vector field \(V\). Let us take \(\eta \in \mathcal{K}(u)\) with curvature \(\kappa \in \mathcal{Z}_C^2(M_0)\). Namely, we have \([\kappa] = e_0\) and \(p_0^*\kappa = \text{d}\eta\). By Theorem 6.14, there is \(\omega_0 \in \mathcal{Z}_C^2(M_0)\) such that

\[
\omega_0 = A_\ast \kappa + \omega_0.
\]

Moreover, we can exploit our freedom in choosing \(\eta\) and \(c_0\) to get

- \(\omega_0 = \omega_0\), if \(e_0 = 0\),
- \(\omega_0 = A_\ast \kappa\), if \(C = 0\).

Indeed, if \(e_0 = 0\), we just take \(\eta\) to be a flat connection, i.e. \(\kappa = 0\). If \(C = 0\), we take \(c_0 = 0\) so that \(A_\ast \neq 0\) (as \([\omega_0^n] \neq 0\) and pick \(\eta \in \mathcal{K}(u)\) with the property that \(\kappa = \frac{1}{A_\ast} \omega_0\), which is equivalent to \(\omega_0 = 0\). We define the function

\[
Q : M_0 \times \mathbb{R} \to \mathbb{R}, \quad Q(\cdot, A)\omega_0^n = (A\kappa + \omega_0)^n, \quad \forall A \in \mathbb{R}.
\]
The function $\mathcal{P} : M_0 \times \mathbb{R} \to \mathbb{R}$ is determined by the properties

- $\mathcal{P}(q, 0) = 0$, $\forall q \in M_0$,
- $\frac{\partial \mathcal{P}}{\partial A}(q, A) = Q(q, A)$, $\forall (q, A) \in M_0 \times \mathbb{R}$.

There holds

$$Q(A) = \int_{M_0} Q(\cdot, A) \omega^0, \quad P(A) = \int_{M_0} \mathcal{P}(\cdot, A)\omega^n.$$  

From the definition, $Q(\cdot, A_s) \equiv 1$ and we define $\mathcal{I} \subset \mathbb{R}$ to be the maximal interval satisfying

- $A_s \in \mathcal{I},$
- $\min_{q \in M_0} Q(q, A) > 0$, $\forall A \in \mathcal{I}$.

Therefore, $\mathcal{P}$ is strictly increasing in the second coordinate on $M_0 \times \mathcal{I}$. This fact will be crucially used in Proposition [7.16]

**Remark 7.9.** If $C = 0$, then $Q(\cdot, A) = (A/A_s)^n$, $P(\cdot, A) = (A/A_s)^{n+1}$, which implies $\mathcal{I} = (0, +\infty)$ or $\mathcal{I} = (-\infty, 0)$ depending on the sign of $A_s$. If $e_0 = 0$, then $Q \equiv 1$ and $\mathcal{P}(\cdot, A) \equiv A$, which implies $\mathcal{I} = \mathbb{R}$.

We now make use of the notion of H-form from Section [3.1]. The H-forms we use below are associated with $\Omega_0$, $\alpha_0 = 0$ and $\sigma_0 = \eta$, the given $S^1$-connection. Namely, we deal with forms of the type

$$\Omega_{H\eta} = \Omega_0 + d(H\eta) \in \Xi_2^0(\Sigma), \quad H : \Sigma \to \mathbb{R}.$$  

**Remark 7.10.** Let us describe the fibres of the map $H \mapsto \Omega_{H\eta}$. To this end, we compute

$$d(H\eta) = d^h H \wedge \eta + Hp_0^*\kappa. \quad (7.6)$$

Here $d^h H$ is the horizontal part of $dH$:

$$d^h H := dH - dH(V)\eta.$$  

Let us suppose now that $d(H\eta) = 0$. From [7.6], we deduce

$$d^h H = 0, \quad Hp_0^*\kappa = 0.$$  

The first relation tells us that $H$ is invariant along curves tangent to $\ker \eta$.

Thus, if $e_0 \neq 0$, then $H$ is constant, as the holonomy in this case is the whole $S^1$. Evaluating the second relation at a point $q \in M_0$, where $\kappa_q \neq 0$, we conclude that $H = 0$.

On the other hand, if $e_0 = 0$, then we can take the quotient $\Pi_k : \Sigma \to \Sigma_k$ by the holonomy of $\eta$, which is a finite subgroup of $S^1$ (see Lemma [4.6]). The bundle $\Sigma_k \to M_0$ is trivial and admits an angular function $\phi_k : \Sigma_k \to S^1$. We define $\phi : \Sigma \to S^1$ as $\phi := \phi_k \circ \Pi_k$. We conclude that for each $H$ with $d^h H = 0$, there exists $\bar{H} : S^1 \to \mathbb{R}$ such that $H = \bar{H} \circ \phi$.

In our setting, the volume of an H-form can be computed in terms of the integration operator associated with the free $S^1$-action $u$:

$$u_* : C^0(\Sigma) \to C^0(M_0), \quad u_*(K)(q) := (p_0)_*(K\eta)(q) = \int_{S^1} K(\Phi^V_t(z)) \, dt, \quad \forall q \in M_0,$$

where $z$ is any point in $p_0^{-1}(q)$ and $(p_0)_*$ denotes the integration along the $p_0$-fibres.
Lemma 7.11. For a function $H : \Sigma \to \mathbb{R}$, there holds
\[ \text{Vol}(H\eta) = \int_{M_0} u_*(\mathcal{P}(p_0, H)) \omega^n_*. \]

Proof. Recalling the definition of Vol from (2.1) and using Fubini’s Theorem, we compute
\[
\text{Vol}(H\eta) = \int_0^1 \int_0^1 \int_{\Sigma} H\eta \wedge \left( p_0^*\omega_0 + rH p_0^*\kappa \right)^n \\
= \int_0^1 \int_0^1 \int_{\Sigma} Q(p_0, rH)(p_0^*\omega_*)^n \\
= \int_{\Sigma} \left( \int_0^1 H Q(p_0, rH) \, dr \right) \eta \wedge (p_0^*\omega_*)^n \\
= \int_{M_0} (p_0)_* (\mathcal{P}(p_0, H)\eta) \wedge \omega_*^n \\
= \int_{M_0} u_*(\mathcal{P}(p_0, H)) \omega_*^n,
\]
where in the second equality, we have used that $H\eta \wedge dH \wedge \eta = 0$.

Remark 7.12. When $e_0 = 0$, defining the Calabi invariant as the volume of $H\eta$
\[ \text{CAL}_\omega(H) := \text{Vol}(H\eta) = \int_{M_0} u_*(H) \omega_*^n, \]
recovers the classical definition of the Calabi invariant for the trivial bundle $M_0 \times S^1 \to M_0$. We say that $H$ is normalised with respect to $\omega_*$ if its Calabi invariant $\text{CAL}_\omega(H)$ vanishes, which is equivalent to requiring that the one-form $H\eta$ is normalised according to Definition 2.5. When $H = h \circ \phi$, as in Remark 7.10, then $H$ is normalised if and only if the integral of $h$ is zero.

Lemma 7.13. If a function $H : \Sigma \to \mathbb{R}$ takes values in $\mathcal{I}$, the $H$-form $\Omega_{H\eta}$ is odd-symplectic.

Proof. To prove the statement, we show that $\Omega_{H\eta}$ is non-degenerate on the hyperplane distribution $\ker \eta$. We compute preliminarily
\[
\Omega_0 + d(H\eta) = \Omega_0 + d^bH \wedge \eta + H p_0^*\kappa = (p_0^*\omega_0 + H p_0^*\kappa) + d^bH \wedge \eta.
\]
Let $z \in \Sigma$ be arbitrary, and set $q := p_0(z)$. We have an inverse $\mathcal{L}_z : T_q M_0 \to (\ker \eta)_z$ for the projection $dp_0 : (\ker \eta)_z \to T_q M_0$, which we can use to pull back the restriction of forms on $\ker \eta$ to the tangent space of $M_0$:
\[
\mathcal{L}_z^* \left( \Omega_0 + d(H\eta) \right)_z \wedge = \left( \mathcal{L}_z^* \Omega_0 + d(H\eta) \right)_z \wedge = \left( \mathcal{L}_z^* (p_0^*\omega_0 + H p_0^*\kappa) \right)_z \wedge = (\omega_0 + H(z)\kappa)_q \wedge = \mathcal{Q}(q, H(z))(\omega_*^n)_q,
\]
where we used $\mathcal{L}_z^*(d^bH \wedge \eta) = 0$. The last form is non-degenerate since $\mathcal{Q}(q, H(z)) > 0$. □
Combining this lemma with the stability property proved in Proposition 3.4, we arrive at the following corollary.

**Corollary 7.14.** Suppose that \( p_0 : \Sigma \to M_0 \) is an oriented \( S^1 \)-bundle with \( e_0 = 0 \) and choose \( \eta \) to be a flat connection for \( p_0 \). Let \( \Omega_0 = p_0^* \omega_0 \) be given, where \( \omega_0 \) is some symplectic form on \( M_0 \). For every sufficiently small \( C^2 \)-neighbourhood \( U \) of \( \Omega_0 \) in \( S_C(\Sigma) \), there exist a \( C^2 \)-neighbourhood \( D \) of \( \text{id}_\Sigma \) in \( \text{Diff}(\Sigma) \) and a \( C^2 \)-neighbourhood \( \mathcal{H} \) of the zero function in the space of normalised functions on \( \Sigma \) with the following property:

\[
\forall \Omega \in \mathcal{U}, \exists \Psi \in \mathcal{D}, \exists H \in \mathcal{H}, \quad \Psi^* \Omega = \Omega_0 + d(H\eta).
\]

**Proof.** Let \( \Omega \in S_C(\Sigma) \) be \( C^2 \)-close to \( \Omega_0 \). By standard elliptic arguments (see for instance [Nic07, Chapter 10]), we have \( \Omega = \Omega_0 + da \) for a \( C^2 \)-small normalised one-form \( \alpha \). We apply Proposition 3.4 with \( \alpha_0 = 0 \), \( \alpha_r \equiv \eta \) and \( \alpha_r = r\alpha \) so that \( \Omega_r = \Omega_0 + rda = \Omega_0 + r(\Omega - \Omega_0) \). Indeed, the hypothesis in Remark 3.5(b) are met, as \( \Omega \) is \( C^0 \)-close to \( \Omega_0 \) and \( \sigma_r \equiv \eta \) is closed. Thus, we have the existence of a diffeomorphism \( \Psi : \Sigma \to \Sigma \) isotopic to \( \text{id}_\Sigma \) and of a normalised function \( H : \Sigma \to \mathbb{R} \) such that \( \Psi^* \Omega = \Omega_0 + d(H\eta) \). More explicitly, from (3.2) and (3.5), we see that we have paths \( \{\Psi_r\} \) and \( \{H_r\} \) with \( \Psi_0 = \text{id}_\Sigma \) and \( H_0 = 0 \) satisfying

\[
X_r \in \ker \eta, \quad \left( i_{X_r} (\Omega_0 + r(\Omega - \Omega_0)) + \alpha \right)|_{\ker \eta} = 0, \quad \dot{H}_r = \alpha(V) \circ \Psi_r,
\]

where \( X_r \) is the time-dependent vector field generating \( \Psi_r \) and \( V \) is the vertical vector field of \( \eta \). From these relations, we see that \( \Psi_r \) is \( C^2 \)-close to \( \text{id}_\Sigma \) and \( H_r \) is \( C^2 \)-small. \( \square \)

### 7.3 The local systolic-diastolic inequality for quasi-autonomous H-forms

We introduce H-forms of special type for which we can prove the systolic-diastolic inequality.

**Definition 7.15.** We say that an H-form \( \Omega_{H\eta} \) is **quasi-autonomous**, if there are pairs \( q_{\text{min}}, q_{\text{max}} \in M_0 \) and \( H_{\text{min}}, H_{\text{max}} : \Sigma \to \mathbb{R} \) such that

1. \( \Omega_{H_{\text{max}} \eta} = \Omega_{H \eta} = \Omega_{H_{\text{min}} \eta} \)
2. \( \min_{\Sigma} H_{\text{min}} = H_{\text{min}}(z), \quad \forall z \in p_0^{-1}(q_{\text{min}}) \), \quad \max_{\Sigma} H_{\text{max}} = H_{\text{max}}(z), \quad \forall z \in p_0^{-1}(q_{\text{max}}) \)

We say that \( H : \Sigma \to \mathbb{R} \) is quasi-autonomous, if the form \( \Omega_{H \eta} \) is quasi-autonomous. Thanks to Remark 7.10 we have \( H_{\text{max}} = H_{\text{max}} \) if \( e_0 \neq 0 \). We recover the standard definition of quasi-autonomous Hamiltonians [HZ11, p. 186] if \( p_0 \) is trivial.

**Proposition 7.16.** Let \( \Omega_* \in Z_C(\Sigma) \) be a Zoll odd-symplectic form with bundle \( p_{\Omega_*} = p_0 \). When \( e_0 = 0 \), the reference connection one-form \( \eta \) for \( p_0 \) is assumed to be flat. If \( H : \Sigma \to \mathbb{R} \) is quasi-autonomous with values in \( \mathcal{I} \) and satisfies \( \mathcal{A}_{\text{min}}(\Omega_{H \eta}) \in \mathcal{I} \), \( \mathcal{A}_{\text{max}}(\Omega_{H \eta}) \in \mathcal{I} \), then

\[
P(A_{\text{min}}(\Omega_{H \eta})) \leq \mathcal{I}o(\Omega_{H \eta}) \leq P(A_{\text{max}}(\Omega_{H \eta})).
\]

Moreover, let us consider the following three conditions:

(a) \( H \) is constant when \( e_0 \neq 0 \), or factors through the map \( \phi \) of Remark 7.10 when \( e_0 = 0 \).

(b) The form \( \Omega_{H \eta} \) is Zoll.

(c) Any of the two inequalities above is an equality.
We have (a) ⇒ (b) and (a) ⇔ (c). The implication (b) ⇒ (c) holds if \( \Omega_{H_0} \) belongs to the neighbourhood \( \mathcal{V} \) given in Proposition 7.4 or if

\[
A_{\text{max}}(\Omega_{H_\eta}) - A_{\text{min}}(\Omega_{H_\eta}) < \inf \left\{ a > 0 \mid a = \langle a(C), [\gamma_r] \rangle, \, \text{for some} \, \, [\gamma_r] \in \pi_1(A_b(\Sigma)) \right\}.
\]

Proof. From (7.6), we see that the oriented \( p_0 \)-fibre over \( q_{\text{min}} \) is a closed characteristic for \( \Omega_{H_\eta} \) with action \( \min H_{\text{min}} \). Therefore,

\[
A_{\text{min}}(\Omega_{H_\eta}) \leq \min H_{\text{min}} \leq H_{\text{min}}.
\]  

(7.7)

Since \( \mathcal{P} \) is increasing in the second coordinate \( A \) when \( A \in \mathcal{I} \), we have by Lemma 7.11

\[
P(\min H_{\text{min}}) = \int_{M_0} u_*(\mathcal{P}(p_0, \min H_{\text{min}})) \omega^a
\]

\[
\leq \int_{M_0} u_*(\mathcal{P}(p_0, H_{\text{min}})) \omega^a
\]

\[
= \mathcal{Q}(\Omega_{H_\eta}).
\]  

(7.8)

In the case of \( e_0 = 0 \), this is true since \( \mathcal{I} = \mathbb{R} \) (see Remark 7.9) and we can assume that \( H_{\text{min}} \) is normalised up to adding a constant. The inequality for \( A_{\text{max}} \) is obtained similarly.

Next, we show the implications between (a), (b), and (c). Assuming (a), we have

\[
\Omega_{H_\eta} = \begin{cases} 
p_0^*(\omega_0 + H\kappa), & \text{if } e_0 \neq 0 
p_0^*\omega, & \text{if } e_0 = 0,
\end{cases}
\]

which is Zoll since the value of \( H \) is in \( \mathcal{I} \). This shows (a) ⇒ (b). Moreover, if \( e_0 \neq 0 \), then \( H = H_{\text{min}} = H_{\text{max}} \) is constant. Therefore, all the inequalities in (7.7) and (7.8) are actually equalities, and (c) follows. If \( e_0 = 0 \), then \( H = \bar{H} \circ \phi \) and the normalised functions \( H_{\text{min}} \) and \( H_{\text{max}} \) are both equal to zero. By the same reason as before, (a) implies (c).

We now show that (c) implies (a). Assume for instance, that \( P(\min(\Omega_{H_\eta})) = \mathcal{Q}(\Omega_{H_\eta}) \).

From inequality (7.8) and the fact that \( A \mapsto \mathcal{P}(\cdot, A) \) is strictly increasing for \( A \in \mathcal{I} \), we have that \( H_{\text{min}}(z) = \min H_{\text{min}} \), for all \( z \in \Sigma \). Therefore, \( H_{\text{min}} \) is constant. If \( e_0 \neq 0 \), then \( H = H_{\text{min}} \) and the conclusion follows. If \( e_0 = 0 \), then \( H = H_{\text{min}} + (H - H_{\text{min}}) \), and therefore, \( H \) factors through the map \( \phi \), as \( H_{\text{min}} \) is constant and \( H - H_{\text{min}} \) factors through \( \phi \) by Remark 7.10 since \( \Omega_{H_\eta} = \Omega_{H_{\text{min}} \eta} \).

Finally, let us assume (b), namely that \( \Omega_{H_\eta} \) is Zoll. If \( \Omega_{H_\eta} \in \mathcal{V} \), we know from Proposition 7.4 (ii) that \( \mathcal{X}(\Omega_{H_\eta}) \subset A(p_0) \). Therefore, \( A_{\text{min}}(\Omega_{H_\eta}) = A_{\text{max}}(\Omega_{H_\eta}) \) due to Proposition 7.4 (iii) and (c) follows. On the other hand, let

\[
\epsilon := A_{\text{max}}(\Omega_{H_\eta}) - A_{\text{min}}(\Omega_{H_\eta}) \geq 0
\]

and suppose that \( \epsilon \) is strictly less than the positive generator of \( \langle a(C), \pi_1(A_b(\Sigma)) \rangle \). Let \( p \) be the bundle associated with the Zoll form \( \Omega_{H_\eta} \), and let \( z_0, z_1 \in \Sigma \) be two points such that \( j_p(z_0), j_p(z_1) \in \mathcal{X}(\Omega_{H_\eta}; p_0) \) and

\[
A_{\Omega_{H_\eta}}(j_p(z_0)) = A_{\text{min}}(\Omega_{H_\eta}), \quad A_{\Omega_{H_\eta}}(j_p(z_1)) = A_{\text{max}}(\Omega_{H_\eta}).
\]
If $s \mapsto z_s$ is any path connecting $z_0$ to $z_1$, we obtain an element $[\gamma_s] \in \pi_1(\Lambda(\Sigma))$ by concatenating a short homotopy for $j_p(z_0)$, the path $s \mapsto j_p(z_s)$, and a reverse short homotopy for $j_p(z_1)$. This loop has the property that

$$\epsilon = \langle a(C), [\gamma_s] \rangle,$$

which implies that $\epsilon = 0$ and condition (c) follows.

**Remark 7.17.** By Proposition 7.14 a quasi-autonomous function that is $C^2$-close to the constant $A_\bullet = A(\Omega)$ satisfies the hypotheses of Proposition 7.10. Moreover, for such a function, conditions (a), (b), and (c) above are equivalent since the corresponding odd-symplectic form belongs to $\mathcal{V}$.

Combining the result we have just proven with the theory of generating functions on closed symplectic manifolds, we can give a proof of Theorem 1.20 namely of the local systolic-diastolic inequality in the case of $e_0 = 0$.

**Proof of Theorem 1.20**

Let $\Omega$ be a Zoll-odd symplectic form such that the associated $S^1$-bundle $p_\Omega$ has vanishing real Euler class. Due to Remark 7.14 we may assume that $p_\Omega = p_0 : \Sigma \to M_0$ and $e_0 = 0$. Let $\omega_s$ be the symplectic form on $M_0$ such that $p_0^* \omega_s = \Omega$. By Lemma 4.6 there exists a trivial oriented $S^1$-bundle $p'_\Omega : \Sigma' \to M_0^1$ and a bundle map $\Pi : \Sigma' \to \Sigma$ such that $p'_0 = \Pi^* p_0$. Therefore, thanks to Proposition 7.3 it is enough to prove Theorem 1.20 for trivial bundles. Let $p_0 : M_0 \times S^1 \to M_0$ with angular function $\phi$. Assume that $\Omega$ is $C^2$-close to $\Omega_\bullet$. Due to Corollary 7.14 there is a diffeomorphism $\Psi : \Sigma \to \Sigma$ such that $\Psi = \Psi_1$, where $\{\Psi_t\}$ is a $C^1$-small isotopy with $\Psi_0 = \text{id}_{M_0 \times S^1}$, and

$$\Psi^* \Omega = \Omega_{H_{d\phi}}$$

for a $C^2$-small normalised Hamiltonian $H : M_0 \times S^1 \to \mathbb{R}$. Thus, $\mathfrak{ol}(\Omega) = 0 = \mathfrak{ol}(\Omega_{H_{d\phi}})$, namely $\text{CAL}_{\omega_s}(H) = 0$. Moreover, $\Omega$ and $\Omega_{H_{d\phi}}$ have the same systole and diastole due to Proposition 6.10 as $[\Psi_t]^*$ preserves the set of short homotopies. The closed characteristics in $\mathcal{X}(\Omega_{H_{d\phi}}; p_0)$ are curves $\gamma \in \Lambda(p_0)$ of the type $\gamma(t) = (q(t), t)$, for some loop $q : S^1 \to M$ with small capping disc $\widehat{q} : D^2 \to M$. Thanks to (6.8), the action of these curves recovers the classical Hamiltonian action

$$A_{H_{d\phi}}(\gamma) = \int_{D^2} \widehat{q}^* \omega_s + \int_{S^1} H(q(t), t) dt.$$

Let $\{\varphi_t\}_{t \in [0,1]}$ be the Hamiltonian isotopy on $(M_0, \omega_s)$ up to time one generated by $H$. The maps $\varphi_t$ are $C^1$-close to the identity, as $H$ is $C^2$-small. From the theory of generating functions on arbitrary symplectic manifolds (see [LM95 Proposition 5.11] and [MS98 Proposition 9.31]), $\{\varphi_t\}_{t \in [0,1]}$ is homotopic with fixed endpoints to a $C^1$-small Hamiltonian isotopy $\{\varphi'_{t}\}_{t \in [0,1]}$ generated by a quasi-autonomous normalised Hamiltonian $H' : M_0 \times S^1 \to \mathbb{R}$, namely $\mathfrak{ol}(\Omega_{H_{d\phi}}) = 0 = \text{CAL}_{\omega_s}(H')$. Adapting an argument in [Sch00], we have an action-preserving bijection

$$\mathcal{X}(\Omega_{H_{d\phi}}; p_0) \cong \text{Fix } \varphi_1 = \text{Fix } \varphi'_1 \cong \mathcal{X}(\Omega_{H'_{d\phi}}; p_0).$$

Therefore, $\Omega_{H_{d\phi}}$ and $\Omega_{H'_{d\phi}}$ have the same systole and diastole. The local systolic-diastolic inequality holds for $\Omega_{H'_{d\phi}}$ due to Proposition 7.10 and hence also for $\Omega_{H_{d\phi}}$ and $\Omega$. \qed
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