Evaluation of Scintillation and Thermally Stimulated Luminescence Properties of Cs₂CdCl₄ Single Crystals

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The optical and scintillation properties of undoped, Sb-doped, and Tl-doped Cs₂CdCl₄ single crystals were investigated. We chose Sb³⁺ and Tl⁺ as the dopants of (ns)²-type electronic configurations. In the photoluminescence and scintillation spectra of undoped Cs₂CdCl₄, three emission peaks were observed. The scintillation decay-time constants were small (1.2 and 13 ns) and light yields were low (several tens of photons/MeV). In contrast, Sb³⁺- and Tl⁺-doped Cs₂CdCl₄ showed s–p transitions of the dopant ions, and the light yield was larger than that of undoped Cs₂CdCl₄ (625 and 1000 photons/MeV, respectively). Tl-doped Cs₂CdCl₄ showed relatively large scintillation decay-time constants (270 ns) because of slow energy transfer from the host to Tl⁺. On the other hand, Sb-doped Cs₂CdCl₄ showed significant afterglow behavior and three thermally stimulated luminescence (TSL) glow peaks (93, 168, and 238 °C) with good linearity with respect to the X-ray dose. In conclusion, the light yield of undoped Cs₂CdCl₄ can be enhanced by doping with Sb³⁺ and Tl⁺; TSL characteristics are achieved by doping Sb³⁺.

1. Introduction

In general, luminescent materials, such as those utilized in radiation detectors, show either of two types of luminescence behavior: immediate luminescence upon irradiation and luminescence upon external stimulation after irradiation. Luminescent materials exhibiting such behavior can be used as scintillators and dosimeters, respectively. Dosimeters store ionizing radiation energy as trapped electrons and holes and exhibit luminescence upon external stimulation by light (optically stimulated luminescence, OSL) or heat (thermally stimulated luminescence, TSL). They are mainly used for personal dose monitoring, and should exhibit suitable sensitivity, linear response to dose, and low fading. On the other hand, scintillators convert the absorbed energy of ionizing radiation immediately into low-energy photons. Scintillation detectors are composed of scintillation materials and a photomultiplier tube (PMT), avalanche photodiode, or Si-PMT, and play important roles in medical imaging,¹ underground resources exploration,² and security inspections.³

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Alkali metal halides are promising scintillation materials for the detection of high-energy photons such as X-rays or gamma rays because of their efficient light yields, high energy resolution, optical transparency, and low manufacturing costs, all of which make them favorable for many practical applications. In fact, NaI:Tl and CsI:Tl exhibit high light yields (~40000 and ~50000 photons/MeV, respectively)(4) originating from $6s^2 \rightarrow 6s^2$ transitions of Tl$^+$ ions, and they have been applied to X-ray and gamma ray detections for many years. However, both types of material have certain disadvantages as scintillators. For instance, NaI:Tl exhibits high hygroscopy in ambient atmosphere, which limits its applications. Its manufacture requires controlled humidity and special handling. CsI:Tl is free from the high hygroscopy exhibited by NaI:Tl; it exhibits a long afterglow attributed to the thermal ionization of trapped electrons (Tl$^0$) followed by radiative recombination with trapped holes [V$_{kA}$ (Tl$^+$)](5) and possesses large decay time constants ($r_1 = 600$ ns and $r_2 = 3.4$ μs) under gamma ray excitation,(6) which leads to pulse pile-up in high-count-rate applications. Although the afterglow can be suppressed by codoping with Eu$^{2+}$, Sm$^{2+}$, Yb$^{2+}$, or Bi$^{3+}$,(7–10) the drawbacks of slow decay components have not yet been completely resolved. Thus, new scintillators having minimal hygroscopy, high light yield, and small decay time constants are required.

Cs$_2$CdCl$_4$ is a transparent nonhygroscopic crystal. The Cs$_2$CdCl$_4$ crystal is tetragonal and crystallizes in a K$_2$NiF$_4$-type structure having the space group I4/mmm with $a = 5.26$ Å and $c = 16.88$ Å, and two equivalent molecules per unit cell.(11) The divalent metal ions (Cd$^{2+}$) are surrounded by a tetragonal-type octahedron, slightly compressed [[Cd–Cl (1) = 2.64 Å and Cd–Cl (2) = 2.51 Å]] along the fourfold axis, which coincides with the crystallographic c-axis.(12) This condition gives rise to a ligand field of D$_{4h}$ symmetry on the divalent site. According to a previous paper,(12) Mn$^{2+}$-doped Cs$_2$CdCl$_4$ crystals exhibit intense 3d$^0$–3d$^n$ transitions of Mn$^{2+}$; detailed spectrum analysis indicated that Mn$^{2+}$ substitutes for Cd$^{2+}$ in the site that occupies the center of the octahedron. This result indicates that Mn$^{2+}$-doped Cs$_2$CdCl$_4$ may exhibit efficient scintillation. However, 3d$^0$–3d$^n$ transitions are not appropriate for scintillator applications because of their very long lifetime (~ms). Alternatively, ns$^2$-type ions (e.g., Tl$^+$, Sb$^{3+}$) exhibit s–p transitions in a crystal field with octahedral symmetry, such as KCl and KI.(13,14) The s–p transitions of ns$^2$-type ions, especially for Tl$^+$, are an attractive mechanism for scintillators because of their efficient light yield(1) and fast decay under UV excitation (~several tens of nanoseconds).(15) In addition to Tl$^+$, Sb$^{3+}$ ions are another possible dopant owing to their similar electron configuration to that of Tl$^+$. According to a previous paper,(16) Sb$^{3+}$ can behave as a luminescence center in Cs$_2$Na$_2$MCl$_6$ ($M =$ Sc, Y, La), the crystal structure of which has octahedral symmetry as inCs$_2$CdCl$_4$. Thus, Sb$^{3+}$ ions are expected to function as luminescence centers in Cs$_2$CdCl$_4$. On the other hand, it has been reported that Sb-doped Sr$_2$SnO$_4$ shows TSL glow peaks due to the s–p transitions of Sb$^{3+}$ after irradiation with UV light.(17) This indicates that Sb-doped Cs$_2$CdCl$_4$ may behave as a dosimeter.

In this study, the photoluminescence (PL) and scintillation properties of undoped, Sb-doped, and Tl-doped Cs$_2$CdCl$_4$ crystals were investigated at room temperature using X-ray-excited luminescence spectra, pulse height spectra, and scintillation decay curves. The dosimeter properties of the Sb-doped Cs$_2$CdCl$_4$ crystal were also investigated using TSL measurements.
2. Experimental Procedure

Single crystals of undoped, Sb-doped, and Tl-doped Cs₂CdCl₄ were grown via the Bridgman–Stockbarger method. Prior to crystal growth, raw powders of CsCl (99.999%, High-Purity Chemical, Japan), CdCl₂ (99.9%, High-Purity Chemical, Japan), and SbCl₃ (99.999%, Wako, Japan) or TlCl (99.999%, Wako, Japan) were thoroughly mixed in a stoichiometric ratio and dried at 423 K for one day to remove water. The crystals were grown under vacuum in sealed quartz ampules. The molar concentrations of Sb or Tl in the melt were 0.1 mol%. Crystals larger than 1 mm were used for optical and scintillation characterizations. Powder X-ray diffraction (XRD) measurements revealed that all the samples had the tetragonal Cs₂CdCl₄ structure.

PL measurements were carried out using a spectrofluorometer (Hitachi, Japan, F-7000) in air at room temperature. The PL spectra of undoped Cs₂CdCl₄ were measured using synchrotron radiation in the vacuum ultraviolet (VUV) region as the excitation source at BL-7B at UVSOR, IMS, Japan. The X-ray-induced radioluminescence (RL) spectra were obtained using an X-ray tube operated at 80 kV and 1.2 mA as the excitation source. The spectra were recorded using a Peltier-cooled charge-coupled device (CCD; Andor, US, DU-920-P) combined with a monochromator (Andor, US, SR163i-UV). The pulse height spectra of the samples were obtained using a photomultiplier tube (Hamamatsu, Japan, R7600U) under 662 keV γ-ray irradiation from ¹³⁷Cs. The scintillation temporal profiles were obtained via a time-correlated single-photon counting method with a pulsed X-ray excitation source. To investigate the dosimeter properties, we measured the TSL glow curve with an original setup consisting of a heater (SAT0983a, Sakaguchi, Japan), a thermostat (SCR-SHQ-A, Sakaguchi, Japan), and a photomultiplier tube (PMT; H11890-210, Hamamatsu, Japan). The heating rate for the TSL measurements was set to 0.5 °C/s, and the sample was heated from 36.5 to 290 °C to measure the glow curve. Furthermore, the TSL spectrum was measured using a CCD-based spectrometer (QE pro, Ocean Optics, US), while the sample was heated by an electric heater (SCR-SHQ-A, Sakaguchi, Japan, E. H. Voc) at a constant temperature. The sample was immediately heated to each temperature after irradiation.

3. Results and Discussion

3.1 Samples

Photographs of the undoped, Sb-doped, and Tl-doped Cs₂CdCl₄ crystals are shown in Fig. 1. The samples had cracks and were not hygroscopic.

3.2 Photoluminescence properties

The room-temperature PL spectra of an undoped Cs₂CdCl₄ crystal with excitations at 58 and 200 nm are shown in Fig. 2(a). Emission peaks at 320 and 600 nm are observed with excitations at 58 and 200 nm, and an emission peak at 430 nm is observed only with excitation at 58 nm. The origins of the emission bands at 320 and 430 nm are unclear at present; the possibilities are
electronic transitions at defects, self-trapped excitons (STEs), and Auger-free luminescence (AFL).

AFL is the radiative transition of an electron from the valence band to the uppermost core level, which can be observed only with excitation energies higher than the binding energy of the core level, e.g., 14 eV for CsCl.\(^{(21)}\) In the case of undoped Cs\(_2\)CdCl\(_4\), the two emission peaks did not

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**Fig. 1.** (Color online) Photograph of undoped, Sb-doped, and Tl-doped Cs\(_2\)CdCl\(_4\) crystals.

**Fig. 2.** (Color online) (a) VUV emission spectra of undoped Cs\(_2\)CdCl\(_4\) crystal, and PL spectra (dotted lines) and PL excitation spectra (solid lines) of (b) Sb-doped and (c) Tl-doped Cs\(_2\)CdCl\(_4\) crystals at room temperature. For undoped Cs\(_2\)CdCl\(_4\), the excitation wavelengths are (1) \(\lambda_{ex} = 200\) nm and (2) \(\lambda_{ex} = 58\) nm. For Cs\(_2\)CdCl\(_4\):0.1%Sb, the excitation and emission wavelengths are (1) \(\lambda_{em} = 516\) nm and (2) \(\lambda_{ex} = 277\) nm. For Cs\(_2\)CdCl\(_4\):0.1%Tl, the excitation and emission wavelengths are (1) \(\lambda_{em} = 350\) nm and (2) \(\lambda_{ex} = 250\) nm.
show a threshold-like behavior in the excitation energy. For this reason, AFL is excluded from consideration as the origin.

The excitation and PL spectra of Sb- and Tl-doped Cs₂CdCl₄ are shown in Figs. 2(b) and 2(c), respectively. In the case of the Cs₂CdCl₄:Sb crystal, the excitation spectrum recorded for λₑₓ = 516 nm consists of five bands: a triplet at 254, 277, and 296 nm, and a doublet at 336 and 350 nm. According to previous studies,(13,16,22) the triplet and doublet bands are ascribed to the ¹T₁u → ¹A₁g and ³T₁u → ¹A₁g transitions of Sb³⁺, respectively. A similar splitting of excitation peaks has been observed for alkali metal halide crystals doped with ns²-type ions such as Pb²⁺,(23–25) Tl⁺,(26) and Sn²⁺(27–29) and has been explained in terms of thermal lattice vibrations.(29,30) The adiabatic potential energy surface (APES) of the excited state is split for general distorted lattice configurations, although it is degenerate in the equilibrium configuration of the ground state. Thus, thermal vibrations, i.e., the dynamical Jahn–Teller effect, cause instantaneous splitting of the excitation bands,(29,30) and the observed splitting in the excitation bands in Fig. 2(b) is attributed to this phenomenon. In the emission spectrum recorded at λₑₓ = 277 nm, a broad peak at 516 nm is observed. According to a previous study of Sb³⁺-doped Cs₂NaMCl₄ (M = Sc, La, Y), (16) this broad peak is attributed to the ³T₁u → ¹A₁g transition of Sb³⁺. In general, ns²-type ions doped into alkali metal halide crystals exhibit two types of emission peaks, the so-called Ax and AT emissions, which arise from two minima on the APES of the ³T₁u relaxed excited state. These minima are due to the Jahn–Teller coupling to the εₕ and τ₂ₕ lattice modes.(15) However, the shapes of these emission peaks excited at the other excitation wavelength (254, 296, 336, and 350 nm) are almost the same. This result indicates that the Jahn–Teller interaction working on the ³T₁u APES does not lead to two emissions (the Ax and AT emissions). A similar case was reported for Sb-doped elpasolite-type crystals such as Cs₂NaMCl₆ (M = Sc, Y, La).(16,31) In Cs₂NaMCl₆, the Sb³⁺ ion substitutes at the M³⁺ ion site with octahedral symmetry. Considering that the PL properties of Cs₂NaMCl₆:Sb³⁺ and Cs₂CdCl₄:Sb³⁺ are similar, the Sb³⁺ ion should substitute at the Cd²⁺ site, which occupies the center of octahedral symmetry as in [SbCl₆]³⁻.

In the emission spectrum of the Cs₂CdCl₄:Tl crystal recorded at λₑₓ = 250 nm, a narrow peak and a broad peak are observed at 350 and 490 nm, respectively. Considering that the wavelengths of the emission peaks are similar to those of CsCl:Tl,(32–35) the narrow peak (350 nm) and the broad peak (490 nm) are ascribed to the ³P₁ → ¹S₀ transition of Tl⁺ and off-center STEs, respectively. Accordingly, the excitation peak recorded at λₑₓ = 350 nm is attributed to the ¹S₀ → ³P₁ transition of Tl⁺.

### 3.3 Scintillation properties

Figure 3 shows the X-ray-induced RL (XRL) spectra of the undoped, Sb-doped, and Tl-doped Cs₂CdCl₄ crystals. In the case of undoped Cs₂CdCl₄, three narrow peaks at 320, 360, and 400 nm, a weak band at 450 nm, and a broad peak with a maximum intensity at 600 nm are observed. In the PL spectra with VUV excitation, the emission peaks at 320 and 360 nm may not be resolved. The emission peaks in the RL spectrum have the same origin as those in the PL spectra.
In the spectrum of the Cs$_2$CdCl$_4$:Sb crystal, a broad peak at 520 nm, which is ascribed to the $^3T_{1u} \rightarrow ^1A_{1g}$ transition of Sb$^{3+}$, is observed. On the other hand, in the spectrum of the Cs$_2$CdCl$_4$:Tl crystal, in addition to the peak at 350 nm, which is also seen in the PL spectrum, we observed a broad intense peak at 600 nm, which is not seen in the PL spectra. This case is similar to that of CsI:Tl. In the PL spectra of CsI:Tl,$^{36}$ there are four peaks; the two weak ones were ascribed to $A_x$ (375 nm) and $A_T$ (401 nm) emissions of Tl$^+$, while the other two intense peaks (486 and 550 nm) were ascribed to off-center configurations of STEs perturbed by Tl$^+$. In the XRL spectra of CsI:Tl,$^{5}$ an intense broad peak at 550 nm was observed and attributed to the thermal ionization of trapped electrons (Tl$^0$) followed by their radiative recombination with trapped holes. Thus, as for Cs$_2$CdCl$_4$, a similar mechanism may result in the difference between the PL spectra and the XRL spectra.

The pulse height spectra of the undoped, Sb-doped, and Tl-doped Cs$_2$CdCl$_4$ crystals are shown in Fig. 4. To estimate the light yields of these crystals, a Gd$_2$SiO$_5$:Ce (GSO) crystal, whose light yield is 10000 photons/MeV, was used for comparison. The peak at channel 630 corresponds to the full-energy peak of GSO. The course gain was 10 times larger for the undoped Cs$_2$CdCl$_4$ and 50 times larger for Cs$_2$CdCl$_4$:Sb and Cs$_2$CdCl$_4$:Tl than for GSO. The full-energy peaks of the undoped Cs$_2$CdCl$_4$ crystal were not clearly visible, so the light yield of undoped Cs$_2$CdCl$_4$ cannot be estimated. On the other hand, shoulders are observed at channels 250 and 180 in the pulse height spectra of the Cs$_2$CdCl$_4$:Sb and Cs$_2$CdCl$_4$:Tl crystals, respectively, although they are not well resolved. Assuming that these shoulders correspond to the full-energy peaks, and considering the quantum efficiencies of the bialkali photocathode of the PMT for GSO (28% for 430 nm), Cs$_2$CdCl$_4$:Sb (10% for 520 nm), and Cs$_2$CdCl$_4$:Tl (2.5% for 600 nm), the light yields of Cs$_2$CdCl$_4$:Sb and Cs$_2$CdCl$_4$:Tl crystals are estimated to be 600 and 1000 photons/MeV, respectively.

The scintillation temporal profiles of the undoped, Sb-doped, and Tl-doped Cs$_2$CdCl$_4$ crystals are shown in Fig. 5. The scintillation decay-time constants were deduced from the sums of two or three exponential decay functions. Table 1 summarizes the decay-time constants and relative intensities. The decay-time constants for the undoped Cs$_2$CdCl$_4$ crystal are 1.2 and
On the other hand, the decay-time constants for the Sb-doped or Tl-doped Cs₂CdCl₄ crystals are composed of a fast component due to host luminescence (0.8 and 11 ns for the Sb-doped sample and 6 ns for the Tl-doped sample) and slow components due to each dopant ion (270 ns for the Sb-doped sample and 200 ns for the Tl-doped sample). According to the intensity ratio, the emissions due to the s–p transitions of the dopant ions do not contribute significantly to the scintillation.
3.4 Dosimeter properties

To investigate the origin of the afterglow of Cs₂CdCl₄:Sb, TSL glow curves after irradiation with X-rays at doses of 1, 10, and 100 Gy are shown in Fig. 6. The Cs₂CdCl₄:Sb crystal shows glow peaks near 91, 166, and 242 °C. The activation energies are derived using the initial rise method(37) on the assumption of first-order kinetics. The intensities at the glow peaks increased proportionally with dose with a negligible change in the glow curve. Therefore, it is appropriate to consider the TSL process as first-order, although first-order kinetics have not been confirmed in this study.(38) The peaks at higher temperatures were analyzed without eliminating the peak at lower temperature because the glow peaks were well separated in temperature. The estimated values are summarized in Table 2. Figure 7 shows the TSL spectra near the glow peaks (93, 168, and 238 °C). The TSL spectra show the same peak with a maximum intensity near 520 nm, which is attributed to the \( ^3T_{1u} \rightarrow ^1A_{1g} \) transition of Sb\(^{3+}\). Thus, it is implied that the afterglow also resulted from the \( ^3T_{1u} \rightarrow ^1A_{1g} \) transition of Sb\(^{3+}\). Figure 8 shows the relationship between the TSL intensity and X-ray dose, namely, the dose response curve. The dose response curves for all the glow peaks have good linearity. This result indicates that the Sb-doped Cs₂CdCl₄ may be used for dosimetry. In addition, the significant TSL may be a reason for the small light yield of Sb-doped Cs₂CdCl₄. Figure 9 shows the TSL glow curves of the undoped Cs₂CdCl₄ crystal after X-ray irradiation from 1 to 100 Gy. A dominant peak is observed at 93 °C with a small shoulder at 170 °C. From the comparison of the glow curves of undoped and Sb-doped Cs₂CdCl₄, the glow peaks at 93 and 168 °C for Sb-doped Cs₂CdCl₄ are attributed to the trapping sites in the Cs₂CdCl₄ host; on the other hand, the glow peak at 238 °C is attributed to Sb\(^{3+}\)-related centers.

Table 2

| Sample          | Peak temperature (°C) | Activation energy (eV) |
|-----------------|-----------------------|------------------------|
| Cs₂CdCl₄:0.1%Sb| 90.5                  | 0.22                   |
|                 | 165.5                 | 0.51                   |
|                 | 242                   | 0.64                   |

Fig. 6. (Color online) TSL glow curves of a Cs₂CdCl₄:0.1%Sb crystal after X-ray irradiation from 1 to 100 Gy.
4. Conclusion

Undoped, Tl-doped, and Sb-doped Cs₂CdCl₄ crystals were synthesized using the Bridgman–Stockbarger method, and their optical, scintillation, and dosimeter properties were evaluated. In the room-temperature PL spectra, the undoped Cs₂CdCl₄ crystal exhibited three types of emission peaks, and the AFL was excluded as the origin of the emission. In the case of a Sb-doped Cs₂CdCl₄ crystal, the splitting of excitation bands due to the dynamical Jahn–Teller effect and the Aₓ emission of Sb³⁺ were observed. These results are similar to those in the case of Sb-doped Cs₂NaMCl₆ (M = Sc, La, Y) and indicate that Sb³⁺ ions substitute for the Cd²⁺ site, which occupies the center in octahedral symmetry. For the Tl-doped Cs₂CdCl₄ crystal, emissions due to ³P₁ → ¹S₀ transitions of Tl⁺ and off-center STEs were observed. Under X-ray
irradiation, the undoped and Sb-doped Cs₂CdCl₄ crystals showed emission peaks similar to those in the PL spectra. On the other hand, the Tl-doped Cs₂CdCl₄ crystal exhibited an additional emission peak, which is attributed to off-center configurations of STEs perturbed by Ti⁺ or the thermal ionization of trapped electrons (Tl₀) followed by radiative recombination with trapped holes. According to the pulse height spectra, the light yields of undoped Cs₂CdCl₄ were slightly enhanced by doping with Sb³⁺ and Tl⁺ (625 and 1000 photons/MeV, respectively). The scintillation decay-time constants for the undoped Cs₂CdCl₄ crystal were small (1.2 and 13 ns). The Sb- or Tl-doped Cs₂CdCl₄ crystals exhibited, in addition to the fast component due to host luminescence, slow components due to the dopant ions as a relatively small contribution (270 ns and 18% for the Sb-doped sample, and 200 ns and 37% for the Tl-doped sample). Furthermore, significant afterglow was observed for the Cs₂CdCl₄:Sb crystal. The TSL glow curves of Cs₂CdCl₄:Sb are composed of three glow peaks at 91, 166, and 242 °C. The TSL is attributed to the 3T₁u → 1A₁g transition of Sb³⁺ according to the TSL spectra. In addition, the dose response curves for all the glow peaks have good linearity.

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