A Spectroscopic Follow-up Program of Very Massive Galaxies at $3 < z < 4$: Confirmation of Spectroscopic Redshifts, and a High Fraction of Powerful AGNs

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Abstract

We present the analysis and results of a spectroscopic follow-up program of a mass-selected sample of six galaxies at $3 < z < 4$ using data from Keck-NIRPS and VLT-Xshooter. We confirm the $z > 3$ redshifts for half of the sample through the detection of strong nebular emission lines, and improve the $\sigma_{\text{phot}}$ accuracy for the remainder of the sample through the combination of photometry and spectra. The modeling of the emission-line-corrected spectral energy distributions (SEDs) adopting improved redshifts confirms the very large stellar masses of the sample ($M_\star \sim 1.5-4 \times 10^{11} M_\odot$) in the first 2 Gyr of cosmic history, with a diverse range in stellar ages, star-formation rates, and dust content. From the analysis of emission-line luminosities and widths, and far-infrared (FIR) fluxes, we confirm that $\gtrsim 80\%$ of the sample are hosts to luminous hidden active galactic nuclei (AGNs), with bolometric luminosities of $\sim 10^{44-46}$ erg s$^{-1}$. We find that the MIPS 24 $\mu$m photometry is largely contaminated by AGN continuum, rendering the SFRs derived using only 24 $\mu$m photometry to be severely overestimated. By including the emission from the AGN in the modeling of the UV-to-FIR SEDs, we confirm that the presence of the AGN does not considerably bias the stellar masses ($< 0.3$ dex at 1 $\sigma$) and that the very large stellar masses of the sample are not significantly biased toward massive galaxies hosting AGNs.

Key words: cosmology: observations – galaxies: evolution – galaxies: formation – galaxies: high-redshift – galaxies: stellar content – infrared: galaxies

1. Introduction

One of the surprising findings in observational studies of galaxies in the early universe is that, in contrast to the bottom-up assembly of dark matter halos inferred from simulations of $\Lambda$CDM cosmologies, in which low-mass halos form early and subsequently grow via continued accretion and merging to form more massive halos at later times, the stellar component of halos appears to have assembled in an anti-hierarchical, top-down manner, with low-mass galaxies assembling most of their mass and forming most of their stars at later times compared to more massive systems (Thomas et al. 2005; Gallazzi et al. 2006; Pérez-González et al. 2008; Wiklind et al. 2008; Marchesini et al. 2009, 2010; Domínguez Sánchez et al. 2011; Muzzin et al. 2013b; Stefanon et al. 2015).

Closely related to this issue is the intriguing finding that the number density of the most massive galaxies (i.e., those with stellar masses $M_\star > 3 \times 10^{11} M_\odot$) seems to evolve slowly from $z \sim 4$ to $z \sim 1.5$ (Pérez-González et al. 2008; Marchesini et al. 2009, 2010; Muzzin et al. 2013b), suggesting that very massive galaxies were already in place at $z \sim 3.5$ and implying that their stellar content was assembled rapidly in the first $\sim 1.5$ Gyr of cosmic history (Caputi et al. 2015; Stefanon et al. 2015). While theoretical models of galaxy formation and evolution have significantly improved in matching observations in recent years, they are still unable to reproduce the observed number density of galaxies at $3 < z < 4$ with log $(M_\star/M_\odot) > 11.5$ (Fontanot et al. 2009; Marchesini et al. 2009, 2010; Ilbert et al. 2013), unless the model-predicted stellar mass functions (SMFs) are convolved to account for the effect of Eddington bias by $\sim 0.3-0.4$ dex (Henriques et al. 2015), which accounts for much larger than typical random uncertainties on $M_\star$ at these redshifts (e.g., Marchesini et al. 2010; Ilbert et al. 2013; Muzzin et al. 2013b). Furthermore, it remains to be determined to what extent the inferred number density measurements at high $z$ are affected by blending of sources in ground-based images.

In recent years, robust evidence for the existence of massive ($M_\star \gtrsim 10^{11} M_\odot$) galaxies at $z > 3$ was enabled by the well-sampled spectral energy distributions (SEDs) and accurate photometric redshifts delivered by near-infrared (NIR) imaging surveys adopting medium-bandwidth filters in the NIR, i.e., NEWFIRM Medium-Band Survey (NMBS; Whitaker et al. 2011, and the FourStar Galaxy Evolution Survey (zFARGBE; Spitler et al. 2014). The wide-field NMBS robustly revealed, a population of “monster” galaxies at $3 < z < 4$ with log $(M_\star/M_\odot) > 11.4$ (Marchesini et al. 2010), whereas the deeper but pencil beam zFARGBE survey extended the study of massive galaxies at $z > 3$ down to the characteristic stellar mass at $z \sim 3-4$, i.e., log$(M_\star/M_\odot) \sim 10.6$ (Spitler et al. 2014; Straatman et al. 2014). Stefanon et al. (2015) used the
UltraVISTA DR1 (Muzzin et al. 2013a) to search for very massive galaxies at $4 < z < 7$, finding a robust candidate of a monster galaxy at $z \sim 5.6$ with $\log(M_*/M_\odot) \sim 11.6$ and quenched star-formation activity.

In contrast to the local universe where massive galaxies are predominantly quiescent and have old stellar populations with red colors (Blanton et al. 2003), the massive galaxy population at $z \sim 3$ is dominated by dusty, star-forming galaxies ($\sim$40–60% of massive population, Marchesini et al. 2010; Spitler et al. 2014), and includes a significant population of already quiescent galaxies ($\sim$30–40%) with ages consistent with the age of the universe at the targeted redshifts. In addition, from the X-ray and radio detections, as well as the ubiquitous very bright fluxes in MIPS 24 $\mu$m (rest-frame 5–6 $\mu$m), it appears that massive galaxies in the early universe commonly host active galactic nuclei (AGNs), both at $z \geq 3$ (Marchesini et al. 2010; Stefanon et al. 2015), as well as at $z < 3$ (Barro et al. 2013; Förster Schreiber et al. 2014).

Despite the well-sampled SEDs delivered by the NMBS and zFOURGE surveys, there is still ambiguity in the photometric redshift solutions for the rest-frame UV faint massive galaxies at $z > 3$. Quantitatively, among the population of very massive galaxies with $\log(M_*/M_\odot) > 11.4$ at $3 < z < 4$ selected from the NMBS by Marchesini et al. (2010), $\sim$50% could have a photometric redshift $z_{\text{phot}} < 3$ solution if their stellar populations were characterized by both an evolved (i.e., stellar age $\geq 1$ Gyr) and very dusty (i.e., $(A_V) \sim 3$ mag) component. This elusive population of high-$z$ old and dusty galaxies, originally proposed in Marchesini et al. (2010) and effectively non-existent at $z < 1$, appears to play a major role among high-$z$ massive galaxies (Marchesini et al. 2014), and has only recently been targeted for detailed investigations (e.g., Bedregal et al. 2017, G. Brammer et al. 2017, in preparation).

Because of the ambiguity in the photometric redshift solutions for as much as half of the population, it is of paramount importance to spectroscopically measure the redshifts of very massive galaxies at $z > 3$. Recently, Marsan et al. (2015) presented the first spectroscopic confirmation of an ultra-massive ($M_* \sim 3 \times 10^{11} M_\odot$) galaxy at $z = 3.35$, dubbed “The Vega Galaxy,” due to the incredible similarity of the integrated SED to an A0V-star such as Vega. The detailed analysis of its UV-to-FIR SED reveals that most of its stars formed at $z > 4$ in a highly dissipative, intense, and short burst of star formation, and that it is transitioning to a post-starburst phase while hosting a powerful AGN.

In this paper, we present the results of a spectroscopic follow-up program of a stellar-mass-complete sample of six galaxies at $3 < z < 4$ originally selected from the NMBS (Marchesini et al. 2010), aimed at spectroscopically confirming their redshifts and further investigating their stellar populations. The brightest galaxy within this sample has been extensively studied and the results were presented in Marsan et al. (2015). We confirm $z > 3$ redshifts for half of the sample through the detection of nebular emission lines, and improve the $z_{\text{phot}}$ accuracy for the remainder of the sample with the combination of photometry and spectra. We present the SED modeling of the combined NMBS photometry and spectroscopy and discuss the properties of the mass complete sample at $z > 3$.

This paper is organized as follows. In Section 2, we describe the target selection, and present the ground-based spectroscopy and data reduction of the mass-selected sample in Section 3. In Section 4, we present the analysis of the spectra, update the redshifts for our targets ($z_{\text{spec}}$, if emission lines are detected; otherwise $z_{\text{cont}}$, obtained through including binned spectra in the modeling of the observed SEDs), the modeling of the SEDs with updated redshifts, and the investigation of the AGNs content of the sample. The results are summarized and discussed in Section 5. We assume $\Omega_M = 0.3$, $\Omega_{\Lambda} = 0.7$, $H_0 = 70 \text{ km s}^{-1} \text{Mpc}^{-1}$, and a Kroupa (2001) initial mass function (IMF) throughout the paper. All magnitudes are in the AB system.

2. Target Selection

The targets presented in this paper are primarily selected from the stellar mass complete sample of massive galaxies log $(M_*/M_\odot) > 11.4$ at $3 < z < 4$ presented in Marchesini et al. (2010), constructed using the NMBS (Whitaker et al. 2011). The NMBS is a moderately wide, moderately deep near-infrared imaging survey (van Dokkum et al. 2009) targeting the COSMOS field (Scoville et al. 2007) and AEGIS strip (Davis et al. 2007). The NMBS photometry presented in Whitaker et al. (2011) includes deep optical $ugriz$ data from the CFHT Legacy Survey, deep Spitzer-IRAC and MIPS imaging (Sanders et al. 2007), Galaxy Evolution Explorer (GALEX) photometry in the FUV (150 nm) and NUV (225 nm) passbands (Martin et al. 2005), NIR imaging with NEWFIRM using the five medium-band filters $J_1$, $J_2$, $J_3$, $H_4$, $H_5$, broadband $K$ (van Dokkum et al. 2009) for both fields. The COSMOS field additionally includes deep Subaru images with the $B_2 V J r i z$+ broadband filters (Capak et al. 2007), Subaru images with 12 optical intermediate-band filters from 427 to 827 nm (Taniguchi et al. 2007), and $IHK_s$ broadband imaging from the WIRCam Deep Survey (McCracken et al. 2010).

The medium-bandwidth NIR filters in NMBS allows the Balmer/4000 Å break of galaxies at $1.5 < z < 3.5$ to be more finely sampled compared to the standard broadband NIR filters providing more accurate photometric redshift estimates (Whitaker et al. 2011). The public NMBS catalog uses photometric redshifts determined with the EAZY code (Brammer et al. 2008) modeling the full FUV-$8 \mu$m SEDs. Stellar mass and other stellar population parameters were determined using Fitting and Assessment of Synthetic Templates (FAST; Kriek et al. 2009), adopting the redshift solution output by EAZY (i.e., $z_{\text{spec}}$), the stellar population synthesis models of Bruzual & Charlot (2003), the Calzetti et al. (2000) reddening law with $A_V = 0$–4 mag, solar metallicity, and an exponentially declining star-formation history (SFH). These SED-modeling assumptions are the same used in the analysis presented in Marchesini et al. (2010).

Table 1 lists the properties of the six targets selected for spectroscopic follow-up. Four objects are the brightest candidates of very massive galaxies at $z > 3$ selected from Marchesini et al. (2010), namely C1-15182, C1-19536, A2-15753, and C1-23152 (i.e., the Vega galaxy, already studied in detail in Marsan et al. 2015). The spectroscopic program targeted two additional galaxies, C1-2127 and C1-19764, consistent with being at $z > 3$ but not originally included in the sample presented in Marchesini et al. (2010) because they did not strictly satisfy the completeness limit in stellar mass (though both massive with log $(M_*/M_\odot) \sim 11.2$–11.3). Table 1 also lists C1-21316 from the sample of Marchesini et al. (2010) but not targeted by our spectroscopic follow-up program, since it has a previously confirmed spectroscopic redshift of $z_{\text{spec}} = 3.971$ (Capak et al. 2010; Smolčić et al. 2012; but see also Miettinen et al. 2015).
### Table 1
Ground-based Spectroscopic Observations

| ID         | α(J2000)     | δ(J2000) | K (mag) | $z_{\text{peak}}$ | log $M^*$ ($M_\odot$) | $\lambda_{\text{range}}$ ($\mu$) | $t_{\text{exp}}$ (minutes) | FWHM (″) | $\lambda_{\text{range}}$ ($\mu$) | $t_{\text{exp}}$ (minutes) | FWHM (″) |
|------------|--------------|----------|---------|-------------------|------------------------|-------------------------------|---------------------------|---------|-----------------------------|---------------------------|---------|
| C1-2127    | $10^\circ00^\prime12^\prime.96$ | $+02^\circ12^\prime11^\prime.5$ | 21.69 $\pm$ 0.07 | 3.17 $\pm$ 0.12 | 11.26 $\pm$ 0.13 | 1.94–2.37 | 60 | 0.5 | ... | ... |
| C1-15182   | $09^\circ59^\prime24^\prime.39$ | $+02^\circ25^\prime36^\prime.5$ | 21.62 $\pm$ 0.09 | 3.56 $\pm$ 0.11 | 11.54 $\pm$ 0.04 | 2.04–2.46 | 60 | 0.5 | 0.3–2.4 | 60 | 0.5 |
| C1-19536   | $09^\circ59^\prime31^\prime.82$ | $+02^\circ30^\prime18^\prime.2$ | 21.65 $\pm$ 0.06 | 3.19 $\pm$ 0.07 | 11.55 $\pm$ 0.03 | 1.94–2.37 | 60 | 0.5 | 0.3–2.4 | 60 | 0.8 |
| C1-19764   | $10^\circ00^\prime21^\prime.12$ | $+02^\circ50^\prime33^\prime.5$ | 21.81 $\pm$ 0.12 | 3.04 $\pm$ 0.17 | 11.22 $\pm$ 0.03 | 1.94–2.37 | 60 | 0.5 | 0.3–2.4 | 60 | 1.1 |
| C1-21316   | $10^\circ00^\prime19^\prime.74$ | $+02^\circ32^\prime04^\prime.3$ | 22.29 $\pm$ 0.16 | 3.68 $\pm$ 0.12 | 11.52 $\pm$ 0.01 | ... | ... | ... | ... | ... |
| C1-23152   | $10^\circ00^\prime27^\prime.81$ | $+02^\circ53^\prime49^\prime.3$ | 20.31 $\pm$ 0.02 | 3.29 $\pm$ 0.06 | 11.42 $\pm$ 0.01 | 1.48–2.37 | 75 | 0.7 | 0.3–2.4 | 60 | 0.5 |
| A2-15753   | $14^\circ18^\prime30^\prime.83$ | $+52^\circ40^\prime24^\prime.6$ | 22.25 $\pm$ 0.06 | 3.14 $\pm$ 0.10 | 11.40 $\pm$ 0.02 | 1.94–2.37 | 210 | 0.6 | ... | ... |

**Notes.** Listed ID’s are from the NMBS catalog (v4.4). “C1-” and “A2-” refer to the COSMOS and AEGIS fields, respectively. The listed redshifts are the best-fit EAZY redshifts ($z_{\text{peak}}$), quoted stellar masses are the best-fit FAST stellar masses used for the selection of the targets (from the NMBS catalog v4.4, Marchesini et al. 2010). $\lambda_{\text{range}}$ is the wavelength range covered by the instrumental setup; $t_{\text{exp}}$ is the on-source exposure time in minutes; FWHM is the average seeing in arcseconds of the observations.

* C1-21316 was not included in the spectroscopic follow-up program, but it has a spectroscopic redshift present in Capak et al. (2010).

* A detailed study of C1-23152 was already presented in Marsan et al. (2015).
C1-21316 is also a submillimeter galaxy listed as AzTEC 5 in the AzTEC millimeter survey catalog in COSMOS (Scott et al. 2008).

3. Observations and Data Reduction

In this section, we briefly describe the ground-based spectroscopic data obtained from Keck-NIRSPEC and VLT-X-shooter as part of follow-up programs aimed at spectroscopically confirming the existence of very massive galaxies at $z > 3$. Along with the observational techniques, we provide a summary of the data reduction. We refer the reader to Geier et al. (2013) and Marsan et al. (2015) for a detailed description of the spectroscopic data reduction. Table 1 summarizes the wavelength coverage, total exposure time, and seeing of the spectroscopic observations. We note that several targets were observed with both NIRSPEC and X-Shooter (C1-15182, C1-19536, C1-19764, and C1-23152).

3.1. NIRSPEC Spectroscopy

We used NIRSPEC (McLean et al. 1998) on the Keck II telescope for H- and K-band spectroscopy of massive $z > 3$ targets primarily in search of their [O III] emission lines. The observations were carried out on the nights of 2011 February 11–13 as part of the NOAO program 2011A-0514 (PI: Marchesini) with a typical seeing of 0.7", which worsened throughout the sequence of observations due to cloudy variable weather. Observations were conducted following an ABA'B' on-source dither pattern. The orientation of the slit was set to include a bright point source when possible to serve as reference when analyzing and combining the two-dimensional rectified frames. Targets were acquired using blind offsets from a nearby bright star. The alignment of the offset star in the slit was checked before each individual 900 s science exposure and corrected when necessary. Before and after each observing sequence, a spectrophotometric standard and an AV0 star was observed for the purpose of correcting for telluric absorption and detector response.

The data reduction for the NIRSPEC observations used a combination of custom IDL scripts and standard IRAF tasks.\(^8\) Bad pixel masks were created by flagging outlier pixels in dark and flat frames. The cosmic rays on the science frames were removed using L.A.Cosmic (van Dokkum 2001). Each spectrum was sky subtracted using an adjacent spectrum with the IDL routines written by G. Becker (2017, private communication). The sky subtracted frames and the sky spectra were rotated and rectified to a linear wavelength scale using a polynomial to interpolate between adjacent pixels. An absolute wavelength dispersion solution was applied with the use of OH skylines. The standard star frames used to correct for atmospheric absorption and detector response were rectified and reduced in the same manner as the science frames. A one-dimensional spectrum was extracted for each telluric standard star (before and after science observations) by summing all the rows (along the spatial direction) with a flux greater than 0.1 times that of the central row. The average of the one-dimensional telluric star spectra was used to correct the two-dimensional science and spectrophotometric star frames for telluric absorption. The one-dimensional spectrum of the spectrophotometric star was extracted in the same manner, and used to create a response function for flux calibration. The position of an on-slit bright source was used to determine the necessary shifts to align target continuum. The two-dimensional rectified and reduced science frames were combined by weighting according to their signal-to-noise ratio (S/N).

3.1.1. VLT-X-shooter

X-shooter is a single-object, medium-resolution echelle spectrograph with simultaneous coverage of wavelength range 0.3–2.5 μm in three arms (UVB, VIS, NIR; D’Odorico et al. 2006). The observations of our targets were carried out in queue mode as part of the ESO program 087.A-0514 (PI: Brammer) in 2011 May, following an ABA'B' on-source dither pattern using the $11'' \times 11''$ and $11'' \times 9$ slits for the UVB and VIS/NIR arms, respectively. This instrumental setup resulted in a spectral resolution of $R = 4200$, 8250, and 4000 for the NIR, VIS, and UVB arms, respectively. For calibration purposes, telluric and spectrophotometric standard stars were observed in the same setup as science observations. We refer the reader to Table 1 for the seeing conditions (FHWM, \(''\)) and exposure times for each object. The data reduction for the X-shooter observations used custom scripts based on the standard X-shooter reduction pipeline (Modigliani et al. 2010). The calibration steps (master darks, order prediction, flat fields, and the two-dimensional maps for later rectification of the spectra) were run with the default parameters in the pipeline (Goldoni 2011). After these five calibration steps, the echelle spectra were dark-subtracted, flat-fielded, and rectified, and the orders stitched (12, 15, and 16 for the UVB, VIS, and NIR arms, respectively). The sky was then subtracted using adjacent exposures. Standard star observations were reduced with the same calibration data as the science frames and used to correct for telluric absorption and detector response. A final spectrum was created by mean stacking all exposures. We refer to Geier et al. (2013) for a detailed description of reduction steps of X-shooter spectra.

3.1.2. Extraction of One-dimensional Spectra

Following Horne (1986), one-dimensional spectra were extracted by summing all adjacent lines (along the spatial direction) using weights corresponding to a Gaussian centered on the central row with a full width at half maximum equal to the slit width used in each observation. To correct for slit losses and obtain an absolute flux calibration, spectroscopic broad/midband fluxes were obtained by integrating over the corresponding filter curves, and a constant scaling was applied to each spectra individually. A binned, lower resolution spectrum with higher S/N was extracted for each spectra using optimal weighting, excluding parts of spectra contaminated by strong sky emission or strong atmospheric absorption. The resulting spectral resolutions of the binned X-shooter spectra were $R \approx 10–45$, 15–40, and 25–50 for the UVB, VIS, and NIR arms, respectively, while the spectral resolutions of the binned NIRSPEC spectra were $R \approx 30–100$ and 30–300 for the $H$ and $K$ bands, respectively.

4. Analysis and Results

In this section, we built and improved upon previously derived and published stellar population parameters from Marchesini et al. (2010) and Whitaker et al. (2011) by...
modeling the emission-line-corrected SEDs constructed from the combination of the UV-to-8 μm NMBS photometry and the binned spectra. We adopted the measured spectroscopic redshift, $z_{\text{spec}}$, when available, or $z_{\text{cont}}$, the improved photometric redshift estimated by modeling the binned spectrum in combination with the to-8 μm NMBS photometry. [O III] nebular emission lines were detected in the two brightest targets, C1-19536 and C1-15182, allowing for the measurement of the spectroscopic redshift. For the remaining targets with continuum detections (C1-2127, A2-15753, and C1-19764), we combined the binned spectrum and the NMBS photometry to model the finely sampled SEDs with EAZY to derive $z_{\text{cont}}$. In addition to the EAZY template set adopted in Whitaker et al. (2011), we also modeled the SEDs using a template set augmented by an "old-and-dusty" template consisting of a 1 Gyr old single stellar population with $A_V = 3$ mag of dust extinction. Marchesini et al. (2010) found that the inclusion of such a template caused approximately half of the massive galaxies population at $z > 3$ to be consistent with a somewhat lower redshift in the range $2 < z < 3$. For a detailed description of the properties of this template and investigation of its effects on the estimated photometric redshifts, we refer the reader to G. Brammer et al. (2017, in preparation).

The results of the detailed analysis of the spectroscopic data of the brightest K-band candidate from Marchesini et al. (2010), namely C1-23152, were presented in Marsan et al. (2015). This paper presented the first spectroscopic confirmation of an ultra-massive galaxy at redshift $z > 3$, along with a detailed investigation of the stellar population and structural properties of a progenitor of local most massive elliptical galaxies when the universe was less than 2 Gyr old, showing its ultra-compact nature, the presence of an obscured powerful AGN, and discussing its evolutionary path to $z = 0$. We refer the reader to Marsan et al. (2015) for the derived stellar population properties of this source.

4.1. Emission-line Features and Spectroscopic Redshift

The extracted one-dimensional spectra were used to measure spectroscopic redshifts ($z_{\text{spec}}$) when nebular emission lines were detected. The nebular emission lines that we set out to measure were primarily Ly$\alpha$ and [O III]λ4959, 5007. The [O III] doublet was identified in the spectra of C1-19536 and C1-15182, along with a weak Ly$\alpha$ detection in the spectrum of C1-19536. A symmetric Gaussian profile plus a local continuum was used to fit these lines in the extracted one-dimensional spectra. Identical redshifts, FWHM, and a ratio of 1:3 for the amplitudes of the [O III]4959, 5007 doublet were assumed. Figure 1 shows the observed one-dimensional spectra around the regions of the considered spectral features for C1-19536 and C1-15182. Red solid curves indicate the best-fit single Gaussian profiles to the emission lines in 1D spectra. The gray shaded regions indicate regions of the spectra significantly affected by telluric sky-lines.
A Monte Carlo approach was used to measure the uncertainties in the centroid, flux, and width of the fitted emission lines. For each 1D spectrum, 1000 simulated spectra were created by perturbing the flux of the true spectrum at each wavelength by a Gaussian random amount with the standard deviation set by the level of the 1σ error at that specific wavelength. Line measurements were obtained from the simulated spectra in the same manner as the actual data. We computed the formal lower and upper confidence limits by integrating the probability distribution of each parameter (centroid, width, continuum, and emission-line flux) from the extremes until the integrated probability is equal to 0.1585. We calculated the fluxes of the observed emission lines by integrating the fitted Gaussian profiles, with uncertainties determined using the 1σ distribution of integrated fluxes from Monte Carlo simulations as described above. The equivalent widths were calculated individually for the observed and simulated spectra by dividing the integrated line flux by the continuum of the fit, with uncertainties derived as above.

The [O III]λλ4959, 5007 doublet emission is detected in both NIRSPEC and Xshooter 1D spectra for C1-15182 and C1-19536, respectively, typical of AGNs of similar luminosity. For C1-21316, the photometric redshift is smaller by ∼6% in Δz/(1 + z) compared to the spectroscopic redshift, but consistent with it at the ∼2σ level. For C1-21316, the photometric redshift is smaller by ∼6% in Δz/(1 + z) compared to the spectroscopic redshift provided by Capak et al. (2010), but still consistent with it at the ∼2σ level. We note, however, that the corresponding spectrum from which zspec was calculated for C1-21316 has been deemed to be of poor quality (Miettinen et al. 2015). The top panel of Figure 2 shows the comparison between the photometric redshifts from Marchesini et al. (2010) and the spectroscopic redshifts or the improved zcont. The [O III] line widths listed in Table 2 are ≈768 ± 152 km s⁻¹ and 871 ± 151 km s⁻¹ for C1-15182 and C1-19536, respectively, typical of AGNs of similar L([O III]) at low-z (Hao et al. 2005) and intermediate-z (Harrison et al. 2016).

### 4.2. Improved Photometric Redshifts

When no emission lines are detected, we used the 1D binned spectra combined with the full FUV-8 μm broad- and medium-band photometry from the NMBS to determine more accurate photometric redshifts (zcont). Following Muzzin et al. (2013a), the default template set used in this work consists of nine templates: the six templates taken from the optimized template set of EAZY, but augmented with emission lines; a template of a 12.5 Gyr old single stellar population constructed using the Maraston (2005) models; a 1 Gyr old post-starburst template; and a slightly dust-reddened Lyman break template. The SEDs of the galaxies in our sample are shown in Figure 3 (left panels), with the EAZY redshift probability functions plotted on the right panels along with the spectroscopic redshift (when available; red), the zpeak from Marchesini et al. (2010; gray), and the zcont (black).

Marchesini et al. (2010) found that up to ∼50% of the massive 3 < z < 4 galaxy sample could be contaminated by a previously unrecognized population of massive, old, and very dusty galaxies at z < 3. As in Marchesini et al. (2010), we considered the case of adding an additional template representing an old (1 Gyr; τ = 100 Myr) and very dusty (A_v = 3 mag) galaxy and repeating the redshift analysis. The resulting redshift probability distributions and best fit zcont are overplotted in the right panels of Figure 3 (dotted–dashed blue). The resulting photometric redshift estimates of only two galaxies in our sample (A2-15753 and C1-21316) are

### Table 2

| Feature      | λobs (Å) | z        | FWHM (km s⁻¹) | EWobs (Å) | L (10⁴₂ erg s⁻¹) |
|--------------|----------|----------|---------------|-----------|-----------------|
| C1-19536     |          |          |               |           |                 |
| Lyα          | 1215.24  | 3.1754   | 1533.4 ± 284.5| 284.0 ± 284.5| 11.5 ± 11.9    |
| O III(2500)  | 1037.68  | 3.3239   | 520.6 ± 142.5 | 142.5 ± 142.5| 26.55 ± 26.55  |
| O III(NIRSPEC)| 1037.68  | 3.3239   | 636.1 ± 124.2 | 124.2 ± 124.2| 17.62 ± 17.62  |
| C1-15182     |          |          |               |           |                 |
| O III(2500)  | 1037.68  | 3.3239   | 548.8 ± 12.2  | 12.2 ± 12.2 | 16.0 ± 16.0    |
| O III(NIRSPEC)| 1037.68  | 3.3239   | 221.7 ± 22.4  | 22.4 ± 22.4 | 7.9 ± 7.9      |

Note. λobs is the laboratory rest-frame wavelength of the targeted spectral lines; z is the derived redshift; FWHM is the intrinsic velocity width of the spectral lines; EWobs is the observed equivalent widths of the spectral lines; and L is the integrated line luminosity calculated using the adopted systemic redshift zspec.
influenced by the addition of the old and dusty template in the analysis, yet both estimates still formally lie at $z > 3$. Only the photometric redshift of C1-19764, not included in the Marchesini et al. (2010) sample, moves below $z = 3$ when the old and dusty template is including in EAZY, with a formal solution of $z_{\text{cont}} = 2.72^{+0.26}_{-0.20}$.

We note that the spectroscopic redshifts of the galaxies in our sample tend to be different from the photometric redshifts at more than $1\sigma$. Specifically, we can see from the right panels of Figure 3 that EAZY overestimates the redshift of two-thirds of our targets by $\sim 0.17$ and $\sim 0.21$ for C1-15182 and C1-19536, corresponding to $\sim 4\%$ and $\sim 5\%$ in $\Delta z/(1 + z)$, respectively.

4.3. SED Modeling and Stellar Population Properties

In order to robustly constrain the stellar population parameters, all spectra and photometry must be corrected for contamination from nebular emission lines. For sources with detected nebular emission lines, we corrected the observed photometry for emission-line contamination by comparing the observed-frame equivalent width of each emission line to the bandwidth of the corresponding filter. We find that the contribution due to [O III] emission is $\sim 15\%$ and $\sim 5\%$ for C1-19536 and C1-15182, respectively, for the NMBS K and $K_s$ filters. Due to the highly uncertain equivalent width calculated for the Ly$\alpha$ emission line of C1-19536, we chose to remove the contaminated filter (Subaru, IA505) when proceeding with the SED fittings.

We estimated the stellar population properties by fitting the binned X-shooter and NIRSPEC spectra in combination with the broadband and medium-band photometry with stellar population synthesis (SPS) models. We used FAST (Fitting and Assessment of Synthetic Templates; Kriek et al. 2009) to model and fit a full grid in metallicity, dust content, age, and star-formation timescale. Stellar population synthesis models of Bruzual & Charlot (2003) were adopted assuming a Kroupa (2001) IMF, an exponentially declining SFH and the Calzetti et al. (2000) extinction law. The age of the stellar population’s fit ranged between 10 Myr and the maximum age of the universe at the redshift of sources with a step size of 0.1 dex. We adopted a grid for $\tau$ between 3 Myr and 10 Gyr in steps of 0.10 dex and allowed the dust attenuation ($A_V$) to range from 0 to 7 mag with a step size of 0.01 mag. We initially modeled the observed SED with the metallicity as a free parameter ($Z = 0.004$, 0.08, 0.02, 0.05) and repeated the modeling by treating the metallicity as a systematic uncertainty by fixing it to solar metallicity ($Z_{\odot}$). Figure 4 shows the observed SEDs of our targets. The results of the SED modeling and the corresponding 1$\sigma$ errors are listed in Table 3. The bottom panel of Figure 2 shows the comparison between the stellar masses from Marchesini et al. (2010) and the stellar masses derived from the improved redshifts (either $z_{\text{spec}}$ or $z_{\text{cont}}$).

As shown in the lower panel of Figure 2, the stellar masses obtained when adopting the improved redshifts and better sampled SEDs are consistent with those based only on the NMBS photometry (Marchesini et al. 2010), spectroscopically confirming the existence of ultra-massive (i.e., log$(M_*/M_{\odot}) > 11.4$) galaxies at $z > 3$, when the universe was younger than 2 Gyr. We note that, although error bars on both axes in the panels of Figure 2 correspond to $1\sigma$ limits, the stellar population parameter space over which the SEDs are modeled are not identical (the observed SEDs are modeled in this work with a finer grid in parameter space than Marchesini et al. 2010, resulting in larger $1\sigma$ errors). For C1-19536 and C1-15182, which have detected nebular emission lines, we find that removing the emission-line contamination to the observed SEDs decreases the stellar masses by $\sim 0.1$ dex (still consistent with being above the mass completeness threshold in Marchesini et al. 2010). The stellar mass of A2-15753 is affected the most when modeling the observed SED including the binned 1D spectra, resulting in a stellar mass of log $(M_*/M_{\odot}) = 11.14^{+0.13}_{-0.11}$, i.e., $\sim 0.25$ dex smaller than what was derived in Marchesini et al. (2010).

The stellar mass of C1-2127 and C1-19764, not part of the sample of ultra-massive galaxies at $z > 3$ from Marchesini et al. (2010), remain effectively the same, i.e., log $(M_*/M_{\odot}) = 11.23^{+0.12}_{-0.14}$ and $11.12^{+0.20}_{-0.22}$, C21316 is the only source with a larger stellar mass when remodelling with improved redshift. This is not surprising because the updated
Figure 3. Left: observed SEDs from the combination of the medium- and broadband NMBS photometry (black filled circles) and the binned spectra (NIRSPEC, green; X-shooter, purple). The gray curves represent the best-fit EAZY templates using the default EAZY template set adopted in this work (light gray when modeling only the medium- and broadband photometry; dark gray when also including the 1D binned spectra). The blue dotted–dashed curves represent the best-fit EAZY template when the additional old and dusty template is included in the estimate of the photometric redshifts. Right: the EAZY redshift probability distributions. The gray curves represent the distribution calculated using only medium- and broadband photometry (NMBS v4.4), whereas the black curves represent the resulting redshift distribution including the 1D binned spectra in EAZY. The blue dotted–dashed curves show the resulting redshift probability distributions when the old and dusty SED is included in the template set of EAZY to estimate the photometric redshifts. Vertical lines indicate the $z_{\text{peak}}$ output by EAZY for different template sets with 1σ uncertainties. Red vertical lines indicate spectroscopic redshifts when available.
The spectroscopic redshift of the source is $\Delta z \sim 0.3$ greater than the $z_{\text{phot}}$ quoted in Marchesini et al. (2010). The time elapsed since the onset of star formation for our targets ranges from $\sim 400$ Myr to $\sim 1.25$ Gyr, consistent with having formation redshifts of $z_{\text{form}} \sim 5-9$. For roughly half of the sample, the best-fit SFR-weighted mean ages $\langle t_{\text{SFR}} \rangle$ (as defined in Föster Scheriber et al. 2004) are $\sim 800$ Myr. Two of the massive galaxies (C1-19764, with $z_{\text{cont}} < 3$ and C1-2127) in our sample have best-fit SFR estimates from SED modeling consistent with low SFRs ($< 1 M_\odot \text{yr}^{-1}$), while the remaining galaxies have SFRs on the order of few tens to hundreds of solar masses per year. We stress that the estimated SFRs, stellar ages, and dust extinction values have large uncertainties.

Figure 5 shows the rest-frame $U - V$ versus $V - J$ colors of the targets using EAZY (Brammer et al. 2008). A Monte Carlo approach was used to measure the uncertainties in the $U - V$ and $V - J$ colors. Specifically, 1000 photometry catalogs were created by perturbing each flux by a Gaussian random number with the standard deviation set by the level of each flux error. The simulated catalogs were each fit with EAZY separately, and the formal upper and lower limits were obtained in a similar manner as for the emission-line fits. The rest-frame colors and 1σ confidence limits for the galaxies in our sample are plotted in Figure 5 in solid colored circles (color scheme identical to Figure 2). Figure 5 also shows color–color evolution tracks of Bruzual & Charlot (2003) models stellar population synthesis models for different SFHs and dust extinction values. The orange solid and dashed tracks represent the evolution of an exponentially declining SFH with $\tau = 100$ Myr and $A_V = 0$ and 2 mag, respectively. The blue solid and dashed tracks represent the evolution for a constant SFH with $A_V = 0$ and 2 mag, respectively. The evolution of an exponentially declining SFH with $\tau = 500$ Myr and $A_V = 1$ is indicated by the dashed green track.

Based on rest-frame colors, only one (C1-19764, only galaxy with best-fit $z_{\text{cont}} < 3$) of the targets falls firmly in the quiescent
Table 3
Best-fit Stellar Population Parameters

| ID      | $z$     | $\log \tau$ (yr) | Metallicity | $\log$(Age) (yr) | $A_V$ (mag) | $\log(M^*) (M_\odot)$ | SFR ($M_\odot$ yr$^{-1}$) | $\chi^2$ |
|---------|---------|------------------|-------------|------------------|-------------|----------------------|-----------------------------|---------|
| C1-2127 | 3.19$^{+0.08}_{-0.07}$ | 6.70$^{+3.10}_{-2.00}$ | 0.000 $^{+0.002}_{-0.003}$ | 7.90$^{+0.82}_{-0.63}$ | 2.00$^{+1.10}_{-0.50}$ | 11.23$^{+0.16}_{-0.12}$ | 0.0067$^{+0.244}_{-0.006}$ | 1.12    |
| C1-15182 | 3.37$^{+0.013}_{-0.011}$ | 8.00$^{+0.02}_{-0.00}$ | 0.02 | 8.60$^{+0.63}_{-0.08}$ | 4.00$^{+0.50}_{-0.18}$ | 11.36$^{+0.19}_{-0.12}$ | 0.03$^{+0.34}_{-0.05}$ | 1.11    |
| C1-19536 | 3.13$^{+0.008}_{-0.005}$ | 8.30$^{+0.14}_{-0.08}$ | 0.05 | 9.00$^{+0.14}_{-0.10}$ | 6.00$^{+0.10}_{-0.15}$ | 11.41$^{+0.05}_{-0.04}$ | 8.20$^{+0.26}_{-0.09}$ | 1.85    |
| C1-21316 | 3.971 | 8.70$^{+0.00}_{-0.00}$ | 0.05 | 9.10$^{+0.20}_{-0.15}$ | 6.00$^{+0.10}_{-0.15}$ | 11.41$^{+0.05}_{-0.04}$ | 9.55$^{+0.04}_{-0.04}$ | 1.35    |
| C1-19764 | 2.72$^{+0.26}_{-0.25}$ | 2.20$^{+0.00}_{-0.00}$ | 0.02 | 9.10$^{+0.20}_{-0.15}$ | 6.00$^{+0.10}_{-0.15}$ | 11.41$^{+0.05}_{-0.04}$ | 8.20$^{+0.26}_{-0.09}$ | 1.85    |
| A2-15753 | 3.1$^{+0.07}_{-0.07}$ | 8.90$^{+0.10}_{-0.10}$ | 0.05 | 9.10$^{+0.20}_{-0.15}$ | 6.00$^{+0.10}_{-0.15}$ | 11.41$^{+0.05}_{-0.04}$ | 4.38$^{+0.14}_{-0.04}$ | 0.81    |

Note. Estimated stellar population parameters from the modeling of the binned UV-to-NIR spectra in combination with the broad- and medium-bandwidth UV-to-8 μm photometry from NMBS derived assuming the Bruzual & Charlot (2003) stellar population synthesis models and an exponentially declining SFR. Spectroscopic redshifts are listed when available, otherwise best-fit EAZY outputs are listed. Quoted errors are 1σ confidence intervals output by FAST and EAZY. The values in parentheses correspond to 3σ confidence intervals. A Kroupa (2001) IMF and a Calzetti et al. (2000) extinction law is assumed in all cases.

4.4. Active Galactic Nuclei Content

4.4.1. Infrared Spectral Energy Distribution

We found that five-sixths of our targets have significant (>3σ) MIPS 24 μm detections (four already studied in Marchesini et al. 2010), consistent with the high fraction of MIPS-detected sources in the sample of IRAC-selected massive galaxies at $z > 3.5$ over GOODS-North (Mancini et al. 2009). At the redshifts of our targets, the observed 24 μm probes the rest-frame $\sim 8$–7 μm. Emission at these wavelengths can arise from warm/hot dust and polycylic aromatic hydrocarbons (Draene & Li 2007), heated by dust-enhanced star formation or AGNs.

Figure 6 shows the observed far-infrared (FIR) SEDs of our targets with significant MIPS 24 μm detections. We include the Herschel PACS 100 and 160 μm photometry from the PACS Evolutionary Probe (PEP) survey DR1 public data release9 (Lutz et al. 2011; Magnelli et al. 2013) and Herschel SPIRE 250, 350, and 500 μm photometry from the Herschel Multi-tiered Extra-galactic Survey (HerMES) DR3/4 public data release10 (Roseboom et al. 2010, 2012; Oliver et al. 2012; Hurley et al. 2017). These fluxes are listed in Table 4. For C1-21316 (a.k.a. AzTEC 5), we also included the de-boottomed photometry from SMA at 890 μm ($S_{890 \mu m} = 93 \pm 1.3$ mJy; Younger et al. 2007), JCMT AzTEC at 1.1 mm ($S_{1.1 \text{mm}} = 6.5 \pm 1.4$ mJy; Scott et al. 2008), ASTE AzTEC at 1.1 mm ($S_{1.1 \text{mm}} = 4.8 \pm 1.1$ mJy; Aretxaga et al. 2011), and JCMT SCUBA2 at 450 and 850 μm ($S_{450 \mu m} = 25.35 \pm 6.04$ mJy and $S_{850 \mu m} = 11.42 \pm 1.38$ mJy; Casey et al. 2013). Most of our sources are not detected in any of the

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9 http://www.mpe.mpg.de/ir/Research/PEP/DR1
10 http://hedam.lam.fr/HerMES/index/dr4
Herschel bands. Only C1-21316, which is a submillimeter galaxy and a radio-loud AGN, is detected in SPIRE and at longer wavelengths; C1-2127 is detected in the blue and green SPIRE bands, but not at 500 μm. No source is detected in PACS despite all being robustly detected in the MIPS 24 μm band. Figure 6 shows the IR SEDs of the targeted sources, along with the 1, 2, and 3σ upper limits. The red dashed SED represents the average SED of Dale & Helou (2002) for each colorspace is indicated in the corresponding region (Kirkpatrick et al. 2015). The (f(AGN)) for each colorspace is indicated in the corresponding region (Kirkpatrick et al. 2015).

First, we derived dust-enshrouded SFRs from the MIPS 24 μm fluxes using the approach presented in Wuyts et al. (2008), which has become one of the most adopted approaches in the literature. Specifically, this method uses the mean logL_{IR,α=1,...,2.5} of the infrared SEDs of star-forming galaxies provided by Dale & Helou (2002) to calculate SFR_{IR}. This method was validated out to z ~ 3.5 for lower mass (i.e., log(M_*/M_☉) < 11) star-forming galaxies using Herschel data (Wuyts et al. 2011; Tomczak et al. 2016). The calculated total 8–1000 μm rest-frame luminosities and the corresponding SFR_{IR} adapted for a Kroupa (2001) IMF from Kennicutt (1998) are listed in Table 5, third and fourth columns. The corresponding errors are also listed, with and without the uncertainties due to random photometric redshift uncertainties when spectroscopic redshifts are not available. As shown in
Table 5, the random uncertainties on $L_{IR}$ and SFR$_{IR}$ as derived using the Wuyts et al. (2008) approach are dominated by the contribution from random photometric redshift errors. We note that a significant scatter of $\sim 0.2$–0.3 dex was found by Wuyts et al. (2011) and Tomczak et al. (2016) between the MIPS 24 $\mu$m derived SFRs and the SFRs derived including robust Herschel PACS detections. Therefore, we stress that the total error budget of $L_{IR}$ and SFR$_{IR}$ listed in Table 5 as derived using the Wuyts et al. (2008) approach is dominated by this scatter and can be as large as a factor of a few. When the Wuyts et al. (2008) approach is adopted, the $L_{IR}$ of our targets range from $\sim 5 \times 10^{12}L_{\odot}$ to $\sim 5 \times 10^{13}L_{\odot}$, typical of Ultra Luminous IR galaxies (ULIRGs) and Hyper Luminous IR galaxies (HLIRGs; Murphy et al. 2011), corresponding to SFR $\sim 500$–5000 $M_{\odot}$ yr$^{-1}$ of obscured star formation. We used the rest-frame 2800 Å luminosity ($\nu L_{2800}$), determined via the best-fit templates in EAZY (using the same methodology for deriving rest-frame colors described in Brammer et al. 2011), as a proxy for $L_{UV}$, reflecting the contribution of unobscured star formation. We used the approach detailed in Bell et al. (2005) to convert $\nu L_{2800}$ to SFR$_{UV}$. The calculated $L_{UV}$ and SFR$_{UV}$ are listed in Table 5.

The SFRs derived by assuming that all the flux at the observed 24 $\mu$m is associated with dust-enshrouded star formation are $\sim 2$–3 orders of magnitude larger than the SFRs estimated from SED modeling (although values for C1-2127 and A2-15753 are consistent within 1σ uncertainties). We can see in Figure 6 that the median pure-SF IR SED template derived from the library of Dale & Helou (2002), and adopted in the Wuyts et al. (2008) approach, is not an adequate fit to the observed detections and limits. For each object, we overplot the median Dale & Helou (2002) template scaled to the observed MIPS 24 $\mu$m detections as dashed red curves. It can be seen that a model in which the IR emission is due only to obscured star formation cannot reproduce the observed IR SEDs. Specifically, the detections and upper limits in Herschel bands put strong constraints on the overall shape of the IR SEDs, rejecting the scenario in which the observed IR properties are due entirely to obscured star-formation activity, and pointing to the presence of a component with warmer dust temperature (arguably originating from the dusty torus of the AGN) not present in the average template from Dale & Helou (2002). This result was already presented in Marsan et al. (2015) for the Vega galaxy at $z = 3.35$, and we now show that it appears to be a common property of ultra-massive galaxies at $z > 3$.

To quantitatively assess the amount of AGN contribution to the IR SEDs, we used the IR template set of Dale et al. (2014), which includes fractional AGN contributions to the mid-infrared (MIR) radiation. We calculated the mean log$(L_{IR})$ for $\alpha = 1$,..., 2.5 from this template set for varying AGN contributions. For visual ease in comparison, we only plot the templates with AGN contributions of 0%, 20%, 40%, 60%, 80%, and 100% scaled to the observed MIPS 24 $\mu$m photometry (dashed blue curves in Figure 6). The Dale et al. (2014) templates that minimize the $\chi^2$ to detections and are consistent with upper limits for each target are plotted as solid blue curves, with $f$(AGN)$_{MIR}$ indicated in each panel. We note that the specified $f$(AGN)$_{MIR}$ refers to the fraction of AGN light to the integrated mid-infrared (i.e., 5–20 $\mu$m) emission. A significant amount of AGN contribution to the MIR SEDs is required by all ultra-massive galaxies at $z > 3$ to explain the IR SEDs. If 3σ upper limits for the Herschel photometry are adopted, we find a lower limit on the fraction of AGNs to the MIR SEDs ranging from 15% to 50% depending on the galaxy. If the 1σ upper limits are adopted, we find a lower limit on the fraction of AGNs to the MIR SEDs ranging from 40% to 75%. Columns 5–8 of Table 5 list the values or upper limits of $L_{IR}$ and SFR$_{IR}$, and the values or lower limits of $f_{\nu}$(AGN) and $f_{\nu}$(AGN) as obtained from adopting the best-fit templates from Dale et al. (2014). The upper limits were derived adopting the 3σ upper limits in the Herschel photometry. We stress that these newly derived $L_{IR}$ and SFR$_{IR}$ should be considered as upper limits, given that we only have lower limits on the AGN contribution to the IR. We note that the three galaxies with X-ray detections, namely A2-15753, C1-15182, and C1-19536 (see Section 4.4.3 below), are undetected in Herschel while being detected at several $\sigma$ in the MIPS 24 $\mu$m band, $f$(AGN)$_{MIR}$ > 40% (75%), 50% (80%), and 15% (60%), respectively, when adopting the 3σ (1σ) upper limits in Herschel. One target, C1-21316, which has $f$(AGN)$_{MIR}$ $\sim 0.35$–0.50 is detected in radio data, indicative of hosting a radio-loud AGN (Section 4.4.4).

We also used the templates from the observed FIR color-based library of Kirkpatrick et al. (2015) to investigate the AGN continuum contribution to the IR SEDs. In this study, libraries of empirical IR SED templates were created from a large sample of high-redshift ULIRGs with rest-frame mid-IR spectroscopy from the Spitzer Space Telescope Infrared Spectrograph (IRS). We display their color diagnostic in the bottom right panel of Figure 6, with the FIR colors of our targets calculated using photometry from Herschel SPIRE and Spitzer MIPS/IRAC plotted in solid colored circles (color scheme identical to Figures 2 and 5), with arrows indicating 2σ limits. Kirkpatrick et al. (2015) used mid-IR features to determine the strength of the AGN contribution to IR SEDs by jointly fitting the mid-IR spectrum of a prototypical starburst.
| ID          | $S_{24}$ (μJy) | $L_{24}^a$ $(10^{12} L_\odot)$ | SFR$_{24}^a$ $(M_\odot$ yr$^{-1}$) | $L_{\text{IR}}^b$ $(10^{12} L_\odot)$ | SFR$_{\text{IR}}^b$ $(M_\odot$ yr$^{-1}$) | $f$ (AGN)$_{\text{bol}}^c$ % | $f$ (AGN)$_{\text{MIR}}^c$ % | $L_{\text{UV}}$ $(10^3 L_\odot)$ | SFR$_{\text{UV}}$ $(M_\odot$ yr$^{-1}$) |
|-------------|----------------|---------------------------------|-------------------------------------|--------------------------------------|-------------------------------------|-------------------------------|-------------------------------|---------------------------------|-------------------------------|
| C1-2127     | 110.6 ± 10.7   | 9.6 ± 0.9 ($^{+1.0}_{-1.5}$)    | 1049 ± 101 ($^{+120}_{-168}$)      | 7.2 ± 0.7 ($^{+0.8}_{-0.8}$)       | 761 ± 291 ($^{+20}_{-194}$)     | 2.7 ± 1.0                     | ~15                           | 5.3                             | 1.7                           |
| C1-15182    | 78.6 ± 7.3     | 10.3 ± 1.0                       | 1119 ± 104                          | <2.6 ± 0.24                        | <259 ± 96                        | >9.2 ± 3.2                     | >50                           | 5.6                             | 1.8                           |
| C1-19536    | 61.6 ± 6.7     | 4.7 ± 0.5                        | 512 ± 56                            | <3.6 ± 0.4                         | <384 ± 148                       | >4 ± 1                        | >15                           | 4.8                             | 1.6                           |
| C1-21316    | 177.1 ± 7.8    | 46.0 ± 1.2                       | 4993 ± 64                           | 9.1 ± 0.4–14.9 ± 0.7               | 900 ± 325–1516 ± 540            | 4 – 12                        | ~35–50                        | 5.4                             | 1.8                           |
| A2-15753    | 165.7 ± 4.2    | 13.3 ± 0.6 ($^{+1.3}_{-1.1}$)    | 1357 ± 34 ($^{+123}_{-127}$)        | <5.3 ± 0.1 ($^{+0.3}_{-0.3}$)      | <533 ± 188 ($^{+12}_{-0.3}$)     | >10 ± 4                       | >40                           | 8.5                             | 2.7                           |

Notes. The Spitzer 24 μm fluxes and derived $L_{\text{IR}}$ and SFR$_{\text{IR}}$ assuming

a the mean Dale & Helou (2002) IR template

b the best-fit mean Dale et al. (2014) IR template in agreement with Herschel 3σ detection limits; these values are strictly upper limits to the $L_{\text{IR}}$ and SFRs allowed by the observed FIR fluxes and limits

c lists the fraction of AGN luminosity contribution to the total bolometric (5–1100 μm) and MIR (5–20 μm) luminosity (values in paranthesis) for the best-fit mean Dale et al. (2014) templates; these values are lower limits constrained by FIR fluxes and limits. For C1-21316, the two values provided in columns 5 to 8 correspond to values derived assuming the two Dale et al. (2014) templates constrained by either the SCUBA 450 μm or the λ > 800 μm fluxes (solid blue curves in Figure 6). The errors listed for $L_{\text{IR}}$ and SFR$_{\text{IR}}$ were computed using just the 24 μm photometric errors (values not in parentheses) and the combination of the 24 μm photometric errors and the photometric redshift errors (values in parentheses).
(M82) and a pure power law for the AGN component to the rest-frame mid-IR spectroscopy. FIR color-based templates were created by stacking the SEDs of galaxies that fall within each \( S_{250}/S_{24} \) (observed) versus \( S_8/S_{3.6} \) (observed) color box \((f_{\text{AGN}})\) of galaxies that fall into each color box is indicated in the lower right panel of Figure 6. The Kirkpatrick et al. (2015) color-based templates consistent with SPIRE, MIPS, and IRAC fluxes and detection limits of our targets are displayed in violet in Figure 6, scaled to the MIPS 24 \( \mu m \) flux of each galaxy. Using the appropriate color-based templates and the listed fraction of \( L_{\text{IR}} \) attributable to star formation for each template (Table 3 in Kirkpatrick et al. 2015), we re-calculated the dust-shrouded SFRs. The SFRs calculated accounting for the AGN contribution to the FIR SEDs this way are broadly consistent with the observed UV-8 \( \mu m \) SED derived SFRs. Specifically, we calculated the dust-shrouded SFRs for C1-15182, C1-19536, and A2-15753 to be \(<100 M_\odot \text{yr}^{-1}, \sim 100-200 M_\odot \text{yr}^{-1} \) for C1-21316, and \( \sim 200 M_\odot \text{yr}^{-1} \) for C1-2127, in line with SED derived SFRs within uncertainties. It should be noted that the Kirkpatrick et al. (2015) library is more representative of our targets because the templates were created using a large sample of U/LIRGs at higher redshifts \((z \sim 0.3-2.8)\) than compared to Dale et al. (2014) templates, which used local galaxies.

To summarize, the investigation of the FIR SEDs of the massive galaxy sample reveals that none of them can be adequately described as purely star-forming systems. In fact, a MIR AGN contribution of at least \( \sim 40\% \), on average, is necessary to explain the marginal detections in the Herschel bands. It is promising that estimates of AGN contribution using the template sets of Dale et al. (2014) and Kirkpatrick et al. (2015) are consistent with each other. Although in both libraries the fractional AGN contribution is calculated over the rest-frame MIR spectrum, the quasar and starburst templates differ. Dale et al. (2014) uses a single quasar template (median PG quasar spectrum of Shi et al. 2013) to linearly mix with a suite of local normal star-forming galaxies spanning a range in \( \alpha \) to characterize different heating levels \((\alpha = 1 \text{ and } 2.5 \text{ for active and quiescent galaxies, respectively})\) over the 5–20 \( \mu m \) wavelength range. In Kirkpatrick et al. (2015), the mid-IR spectrum of the prototypical starburst M82 is used to represent the star-formation component, and the AGN component is characterized by a pure (free-slope) power law.

4.4.2. [O III] Luminosities and Bolometric Correction

We use the luminosity-dependent O[III] bolometric correction factor, \( C_{\text{OIII}} \) from Lamastra et al. (2009) for the two galaxies with [O III] emission-line detections (namely C1-19536 and C1-15182). \( L_{\text{bol}} \approx 454 \times L_{\text{O III}} \). Because both C1-19536 and C1-15182 have non-negligible best-fit \( A_V \) values from SED fitting, we assume that the observed [O III] line luminosities are a lower limit to the intrinsic luminosity of [O III]. Assuming all the observed [O III] emission is due to AGN radiation, we find \( L_{\text{bol}} > 8 \times 10^{45} \text{ erg s}^{-1} \) and \( >3.5 \times 10^{45} \text{ erg s}^{-1} \) for C1-19536 and C1-15182, respectively, consistent with them hosting powerful hidden AGNs.

4.4.3. Continuum Emission from AGNs

The stellar population parameters derived in Section 4.3 may be influenced by the presence of strong AGN continuum contamination to the observed photometry, most significantly biasing the stellar mass estimates. We investigated this by using the publicly available fully Bayesian Markov Chain Monte Carlo SED fitting algorithm of active galaxies, AGNFitter11 (Calistro Rivera et al. 2016), to model the observed SEDs of the 24 \( \mu m \) detected targets. AGNFitter simultaneously fits the total active galaxy SED by decomposing it into four physically motivated components: the stellar population of the host galaxy, cold dust emission from star-forming regions, the accretion disk emission (Big Blue Bump, BBB) and hot dust emission surrounding the accretion disk (torus). We used the fiducial library of templates that are supplied with AGNFITTER, which has been tested to perform well in classifying Type1 and Type 2 AGNs purely based on observed broadband photometry. We briefly describe the modeling assumptions (as adopted in Calistro Rivera et al. 2016) and the fiducial library of templates used in AGNFITTER, and highlight important distinctions compared to the modeling assumptions employed in previous sections.

The BBB describing the thermal radiation emitted from the accretion disk surrounding the central supermassive black hole was modeled using a single modified template based on the composite spectrum of Sloan Digital Sky Survey Type 1 QSOs (Richards et al. 2006). Extinction to the emitted BBB spectrum was modeled assuming a Small Magellanic Cloud reddening law (Prevot et al. 1984). Emission from the dusty torus was modeled by using a library of SEDs with varying hydrogen column densities \((N_H = 21–25)\) created based on the set of empirical templates of Silva et al. (2004). The BC03 stellar population synthesis models are used to construct the library of templates to describe the host galaxy’s stellar emission. Although the stellar population synthesis model assumptions are similar to those adopted in FAST as presented in Section 4.3, the explored parameter space is somewhat reduced. Specifically, a similar exponential declining SFH is adopted in AGNFITTER, of the form \( SFR(t) \propto e^{-t/\tau}, \) with a mildly more restricted range for the timescale \( \tau \) from 0.1 to 10 Gyr. A model with constant SFR is also included in AGNFITTER. The grid of stellar population ages is created from 0.2 Gyr to the age of the universe at the target’s redshift in step of \( \Delta t \sim 0.1 \text{ dex at fixed solar metallicity. AGNFITTER adopts a Chabrier (2003) and the Calzetti et al. (2000) reddening law to model the extinction to the stellar light. The library of templates for modeling the emission from cold dust in star-forming regions includes the template set of Dale & Helou (2002) (also used in Section 4.4.1) and Chary & Elbaz (2001). For a more detailed description of the fitting algorithm, modeling assumptions, and template libraries, we refer the reader to Calistro Rivera et al. (2016).
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Table 6

| ID      | log M* | Δ log M* | f_{AGN} (0.1–1 μm) | f_{AGN} (1–30 μm) |
|---------|--------|----------|---------------------|-------------------|
| C1-2127 | 11.24  | -0.11    | 0.37               | 0.09              |
| C1-15182| 11.61  | -0.05    | 0.09               | 0.96              |
| C1-15936| 11.37  | -0.08    | 0.10               | 0.98              |
| C1-21316| 11.34  | -0.24    | 0.28               | 0.39              |
| A2-15753| 11.21  | +0.04    | 0.04               | 0.99              |

Note. Results obtained from the modeling of the binned UV–NIR spectra, NMBS UV-8 μm photometry, and FIR detections and upper limits with AGNfitter. The best-fit stellar masses were adjusted to the Kroupa (2001) IMF in order to compare with results listed in Table 3. ΔM* lists the effect on the derived stellar mass when including AGN templates in the modeling of the panchromatic SEDs (Δ log M* = log (M*)_{AGNfitter} − log (M*)_{FAST}) at fixed metallicity (Z = Z⊙). f_{AGN} is calculated as the relative contribution of the AGN templates in the considered wavelengths (i.e., the Big Blue Bump emission describing the accretion disk at 0.1–1 μm and the hot dusty torus emission at 1–30 μm).

Figure 7. SEDs output when modeling the observed UV-FIR photometry of the 24 μm detected targets with four components using AGNfitter (Calistro Rivera et al. 2016). The orange and green curves correspond to the stellar emission and starburst components, respectively. The accretion disk and the surrounding hot dusty torus components are represented as blue and purple curves, respectively. In each panel, the SED components corresponding to 10 randomly selected realizations from the posterior probability density functions are overplotted in order to visualize the range of parameter space. The total SEDs from the realizations are plotted in red.

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Table 6

| ID      | log M* | Δ log M* | f_{AGN} (0.1–1 μm) | f_{AGN} (1–30 μm) |
|---------|--------|----------|---------------------|-------------------|
| C1-2127 | 11.24  | -0.11    | 0.37               | 0.09              |
| C1-15182| 11.61  | -0.05    | 0.09               | 0.96              |
| C1-15936| 11.37  | -0.08    | 0.10               | 0.98              |
| C1-21316| 11.34  | -0.24    | 0.28               | 0.39              |
| A2-15753| 11.21  | +0.04    | 0.04               | 0.99              |

Note. Results obtained from the modeling of the binned UV–NIR spectra, NMBS UV-8 μm photometry, and FIR detections and upper limits with AGNfitter. The best-fit stellar masses were adjusted to the Kroupa (2001) IMF in order to compare with results listed in Table 3. ΔM* lists the effect on the derived stellar mass when including AGN templates in the modeling of the panchromatic SEDs (Δ log M* = log (M*)_{AGNfitter} − log (M*)_{FAST}) at fixed metallicity (Z = Z⊙). f_{AGN} is calculated as the relative contribution of the AGN templates in the considered wavelengths (i.e., the Big Blue Bump emission describing the accretion disk at 0.1–1 μm and the hot dusty torus emission at 1–30 μm).

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larger (by ~0.04–0.05 dex) than when no AGN contamination is assumed. This appears to be caused by slightly older best-fit ages of the stellar population (and hence slightly larger mass-to-light ratios) when the SEDs are modeled accounting for the AGN contribution. We note, however, that these differences are smaller than (or comparable to) the 1σ error on the stellar mass. For the other three galaxies, the stellar masses derived by accounting for the AGN contamination are found to be smaller (by ~0.1–0.3 dex on average) than the stellar masses derived when no AGN contamination is assumed. The best-fit stellar masses derived assuming AGN contamination ranges from 1.5 x 10^{11} M⊙ to 4 x 10^{11} M⊙. All galaxies except C1-21316 are found to be more massive than 10^{11} M⊙ when modeled by AGNfitter. For all galaxies but C1-21316, the systematic effect to the derived stellar masses from AGN contamination is found to be smaller than a factor of ~2 at 1σ. This is because the rest-frame UV–optical SEDs of these galaxies are dominated by the stellar light despite the rest-frame FIR SEDs being dominated by the obscured AGN. For C1-21316, the stellar mass is found to be overestimated by a factor of ~2 on average, although the 1σ uncertainty allows for solutions that are smaller by as much as a factor of ~8. Finally, we note that the stellar masses of all studied galaxies derived from AGNfitter are consistent with the stellar masses derived using FAST within the 1σ uncertainties.

4.4.4. X-Ray and Radio Detections

As previously done in Marchesini et al. (2010), we used the publicly available Chandra X-ray catalogs (Laird et al. 2009 and Elvis et al. 2009 for the AEGIS and COSMOS fields, respectively)
and the improved redshifts to revise the X-ray luminosities of the three detected sources (C1-15182, C1-19536, and A2-15753). Assuming a power-law photon index $\Gamma = 1.9$ (Nandra & Pounds 1994), the X-ray luminosities $L_{2-10\,\text{keV}}$ are $(6.4 \pm 1.7) \times 10^{35}$ for C1-15182, $(16.6 \pm 2.6) \times 10^{35}$ for C1-19536 and $(9.6 \pm 1.5) \times 10^{35}$ erg s$^{-1}$ for A2-15753 (redshift uncertainties are included in quoted errors). We used the empirically derived observed $L_{\text{[O III]}-L_{2-10\,\text{keV}}$ relation of local AGNs from Ueda et al. (2015) to investigate the expected X-ray flux of the sources with detected [O III] emission (C1-19536 and C1-15182). The estimated $L_{2-10\,\text{keV}}$ for these galaxies using the observed $L_{\text{[O III]}}$ are an order of magnitude less than the luminosities calculated using X-ray detections (using $\Gamma = 1.4$ lowers this factor to $\sim 6$–8). We calculated upper limits to the rest-frame $L_{2-10\,\text{keV}}$ for the targets not detected (C1-2127, C1-19764, and C1-21316) using the 3$\sigma$ detection limits with $\Gamma = 1.9$. The 3$\sigma$ upper limits are in the range of $L_{2-10\,\text{keV}} \sim 5$–20 $\times 10^{44}$ erg s$^{-1}$, comparable to only the very powerful X-ray AGN, and therefore we cannot conclude or rule out whether these galaxies harbor AGNs.

From the sample of galaxies presented here, only C1-21316 (AzTEC-5) is detected at 1.4 GHz (Schrinner et al. 2010), with a flux of 0.126 $\pm$ 0.015 mJy. We calculated the rest-frame flux densities assuming the canonical value of the radio spectral index of $\alpha = 0.8$ (Condon 1992). The calculated rest-frame flux density is $L_{1.4\,\text{GHz}} = 10.6 \pm 1.3 \times 10^{27}$ W Hz$^{-1}$, above the threshold to be classified as a radio-loud AGN (log($L_{1.4\,\text{GHz}}$) $> 25$; Schinnerer et al. 2007). The radio spectral index is not observed to evolve significantly over cosmic time (Magnelli et al. 2015), and is consistent with values derived for different samples of galaxies (Ibar et al. 2009, 2010; Ivison et al. 2010; Bourne et al. 2011).

4.4.5. Redshift Evolution of the AGN Fraction

The fraction of AGNs ($f_{\text{AGN}}$) within a galaxy population represents a probe of the level of AGN activity and of the growth of supermassive black holes at the center of galaxies. The energetic output from AGNs is considered to be one of the dominant feedback processes that can regulate, and even halt altogether, the infalling of gas and star formation in massive galaxies. The measurement of the level of AGN activity in massive galaxies as a function of redshift is therefore an indirect probe of the impact of AGN feedback at a given cosmic time.

Figure 8 shows the evolution of $f_{\text{AGN}}$ as a function of redshift for galaxies with stellar mass $M_* > 10^{11} M_\odot$. In our sample of very massive galaxies at $3 < z < 4$, AGNs appear almost ubiquitous, with a very large AGN fraction of $f_{\text{AGN}} \approx 0.8 \pm 0.2$ (in which we assumed that C1-2127 does not host an AGN). Figure 8 also shows the fraction of AGN in massive galaxies at $0.6 < z < 3.0$ from Cowley et al. (2016), at $1.7 < z < 2.7$ from Kriek et al. (2007), and in the local universe from Kauffmann et al. (2003). Our measurement is consistent with the very high AGN fraction of $f_{\text{AGN}} = 0.86^{+0.14}_{-0.22}$ at $2.6 < z < 3.0$ as measured by Cowley et al. (2016) using the zFOURGE data set. On the contrary, the AGN fraction in massive galaxies at $z \sim 0$ is much smaller, $f_{\text{AGN}} \approx 0.3$, and it appears to remain relatively constant all the way out to $z \sim 2.5$. Figure 8 clearly shows a dramatic transition in the fraction of AGNs in galaxies with $M_* > 10^{11} M_\odot$ happening at $z \sim 2.5$, i.e., when the universe was $\sim 3$ Gyr old. We quantified the significance of the transition in the AGN fraction at $z \sim 2.5$ by assuming that the fraction of AGN is constant with redshift, and equal to the bi-weight mean of the observed fraction at $z < 2.5$ (i.e., $f_{\text{AGN}} \approx 0.29$) and testing this hypothesis with the chi-squared statistics of the two points at $z > 2.5$. If each point at $z > 2.5$ is considered by itself, we find a probability $p < 0.0006$ of obtaining each $z > 2.5$ measurement at least this discrepant from the no evolution assumption. If the two $z > 2.5$ measurements of the AGN fraction are considered together, we find a probability $p < 10^{-7}$ of obtaining both $z > 2.5$ points at least this discrepant from the no evolution assumption. At earlier times, most (all) supermassive black holes at the center of massive galaxies appear to be actively accreting, whereas at later times most ($\sim 2/3$) of them are found dormant. Although the small sample does not allow us to robustly quantify a potential enhancement of the AGN fraction in our spectroscopically confirmed sample of very massive galaxies at $z > 3$ due to [O III] emission-line contamination to the stellar continuum. However, we note that the [O III] line emissions detected in the spectra are found to contribute at most 15%, and more generally $\lesssim 5\%$, to the observed $K$-band fluxes. The limited observed [O III] emission-line contamination suggests that our sample is not significantly biased toward strong [O III] emitters, i.e., toward massive galaxies hosting AGNs, and that the derived AGN fraction in our sample is not considerably enhanced. Future wide area surveys will be able to constrain more robustly the fraction of AGN in massive galaxies at $z > 2.5$, whereas deeper surveys will be able to investigate the evolution of the AGN fraction in lower-mass galaxies.

5. Summary and Discussion

In this paper, we investigated the observed-frame UV and NIR spectra of a sample of six $3 \lesssim z_{\text{spec}} < 4$ massive galaxies selected from the NMBS. We confirmed the spectroscopic redshift of two galaxies (C1-15182, $z_{\text{spec}} = 3.371$; C1-19536, $z_{\text{spec}} = 3.139$) through nebular emission lines (Figure 1). We re-modeled the medium- and broadband photometry in combination with the binned one-dimensional spectra of detected sources to derive improved photometric redshifts with the EAZY code ($z_{\text{cont}}$). The best-fit $z_{\text{cont}}$ of only one of our
targets (C1-19764, not included in Marchesini et al. 2010 sample to calculate the high-mass end SMF at $3 \leq z_{\text{phot}} < 4$) is $z < 3$ (although the redshift solution extends to $z > 3$).

We used FAST in conjunction with the BC03 stellar population synthesis models, the Calzetti et al. (2000) extinction law, the Kroupa (2001) IMF, and an exponentially declining SFH to model the observed SEDs (Figure 4, Table 3). We included an additional source in our analysis from the Marchesini et al. (2010) sample, C1-21316, which has a spectroscopic redshift of $z_{\text{spec}} = 3.971$. We find that the updated stellar population parameters are consistent with those previously derived using only medium- and broadband photometry. The stellar masses of galaxies with detected optical emission lines decrease by $\sim 0.1$ dex when emission-line contamination to the observed photometry is accounted for. From the SED modeling, our sample of massive galaxies show ranges in stellar population properties, in accordance with previous studies at $3 \leq z < 4$ (Muzzin et al. 2013b; Marchesini et al. 2014; Spitler et al. 2014). We find that $\sim 50\%$ of sources in this sample are characterized by average stellar ages of $\sim 800$ Myr, suggesting that the bulk of the stellar mass in these systems was formed early in the universe ($\zeta_{\text{form}} \sim 4.5$) (these also have ongoing star formation, with SFR $\sim \text{few} \times 10^{-10} M_{\odot}$ yr$^{-1}$, but these values are highly dependent on model parameters).

All but one of the sources in this sample have significant MIPS 24 $\mu$m detections. The total IR luminosities estimated from the observed 24 $\mu$m assuming IR templates representative of starburst galaxies are $\sim 2$--3 orders of magnitude greater than the SFRs estimated from SED modeling. Inspecting the observed FIR SEDs including Herschel-PACS and SPIRE detections and upper limits effectively rules out dust-shrouded star formation as the only source of the observed 24 $\mu$m flux. In fact, using two separate IR libraries that include different levels of contribution from an AGN (Dale et al. 2014; Kirkpatrick et al. 2015) reveals that $> 15\%$ of the MIR emission must originate from obscured AGNs. However, despite the significant AGN emission contribution to the MIR SEDs, we find that the rest-frame UV–optical light is largely dominated by stellar emission. For all but one galaxy, the presence of the AGN does not considerably bias the stellar mass estimates, with typical systematic differences within $\sim 0.1$ dex, and stellar masses at most a factor of $\sim 2$ smaller at the 1$\sigma$ level. Even after accounting for the presence of the (mostly obscured) AGN, the inferred stellar masses of all but one galaxy are found to be larger than $10^{11} M_{\odot}$ within the 1$\sigma$ uncertainties. The large fraction ($> 80\%$) of massive galaxies hosting AGNs at $3 < z < 4$ inferred from strong [O III] emission lines, IR SED properties, and X-ray and radio detections implies that $3 < z < 4$ must be an extremely active time in the cosmic history for the growth of the supermassive black holes at the massive end of the galaxy population.

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