Charmonium Production at Tevatron, HERA and LHC

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Charmonia hadro- and photo-production are analyzed in the framework of the color-octet model taking into account higher-order effects induced by initial-state radiation of gluons. We argue that color-octet matrix elements obtained from Tevatron data might be reconciled with those extracted from HERA data on inelastic $J/\psi$ photoproduction. Finally we estimate cross sections for events with muons from prompt $J/\psi$’s as a background for the $B$ physics programme at LHC experiments.

1. INTRODUCTION

Over the last years, charmonium physics has recovered a leading role in the challenge of understanding both perturbative and non-perturbative aspects of the strong interaction dynamics. The surplus of $J/\psi$ and $\psi'$ resonance production found at the Fermilab Tevatron triggered an intense theoretical activity as conventional wisdom (i.e. the color-singlet model) failed dramatically to reproduce experimental data. Such a discrepancy, by more than an order of magnitude, led to the introduction by Braaten and Fleming of a new production mechanism, the so-called color-octet model (COM), later cast into a well defined framework based on an effective non-relativistic field theory (NRQCD). However, although those non-perturbative parameters introduced in the model and needed to bring accordance with experimental data can be rigorously defined in the effective theory (as vacuum expectation values of effective field operators), the possibility of obtaining numerical values from lattice calculations seems unfortunately still far. This current theoretical inability means in practice that such long-distance parameters have to be determined from fits to experimental data.

On the other hand, there is the founded conviction that due to the distinct energy scales, charmonium production cross section at high transverse momentum factorizes into a short- and a long-distance part, the former calculable from perturbative QCD, the latter parameterized by NRQCD matrix elements expected to be universal - i.e. the same for different production mechanisms and energies. However, numerical values of some non-perturbative parameters extracted from Tevatron and HERA apparently are in disagreement each other, casting some doubts on the validity of the COM itself, or at least on the above-mentioned universality hypothesis.

This paper is partially devoted to argue that the claim on the inconsistency of the NRQCD matrix elements obtained from hadro- and photoproduction may be premature because the effective intrinsic transverse momentum of partons has been overlooked.

Indeed, it was well known long time ago that particle production in hadron-hadron collisions requires the introduction, usually as a phenomenological input, of a primordial $k_t$ of partons in addition to their expected Fermi motion. During this Conference, we have heard Gemme and Lazzeroni discuss the need of a large intrinsic $k_t$ to reproduce experimental data on $D$ and $B$ production. In fact, most of the effective transverse momentum of partons should be attributed...
to multi-gluon emission by the initial state. In
the literature sometimes this effect is simulated or
simply parameterized by means of a convolution
of the hard process cross section with a Gaussian
$k_t$ distribution whose width is an adjustable pa-
rameter. However, strictly speaking the effect
is modeled more reliably by using a Monte Carlo
generation of the real complexity of the initial-
state radiation. As in our previous analysis of
charmonium hadro-production we have used
PYTHIA managing multi-gluon emission ac-
cording to an elaborated algorithm.

2. EFFECT OF THE INITIAL RADIA-
tion on the $p_t$ DISTRIBUTION

Let us point out that, technically, the inclusion
of multi-gluon emission by itself hardly changes
the overall cross section for charmonium pro-
duction obtained from PYTHIA. However, the
smearing caused by the initial-state radiation is
responsible for a migration of events towards the
high-$p_t$ tail of the charmonium distribution, thus
enhancing the differential cross section in the re-
region where experimental points actually lie. In
figure 1 (see Ref. [11]) this effect is shown, clearly
pointing out the need for a correction of the re-
sults from Cho and Leibovich if a more realistic
description of the underlying hadronic complex-
ity of the interaction (i.e. initial-state radiation)
is required.

In taking into account this $p_t$-dependent $K$
factor, new (lower) values for the long-distance
parameters, and ultimately the color-octet matrix
elements, are obtained from charmonium hadroproduction. In Table 1 we present our val-
ues determined from older Tevatron data together with the new ones on $J/\psi$
production, both showing a better agreement with HERA results than before.

3. THEORETICAL UNCERTAINTIES

There is no space to review all the theoretical
uncertainties related to the extraction of matrix
elements from charmonium production which on
the other hand have extensively been reviewed in
the literature (see for example [4]). Let us only
mention:

- Ambiguities in the choice of energy scales, charm mass, PDF, ...
- Partial breaking of the heavy-quark spin symmetry in charmonia systems, ...
- Subleading $O(\geq v^7)$ contributions, higher-order $\alpha_s$ effects, ...
- Possible kinematic effects near boundaries of phase space
- Possible effects due to the surrounding medium breaking universality

Below we comment in some detail on the possible effects due to the sensitivity to the kinematics
of soft gluon radiation in the long-distance evolution of the $c\bar{c}$ pair into final charmonium.
3.1. Finite width effects

Finite width effects can have consequences as Ernström, Lönnblad and Vänttinen discussed some time ago \[19\]. Indeed, even a small energy difference (of order $M^2$) between the mass of final charmonium and the invariant mass of the created $c\bar{c}$ pair should have an effect on the (steep) momentum distribution. Nevertheless, the point is that the formation of an intermediate colored $c\bar{c}$ state does not preclude the possibility of an energy deficit and not necessarily an excess w.r.t. charmonium mass, the latter implying the emission of soft gluons as usually considered. In other words, the $c\bar{c}$ pair could get energy from the surrounding medium.

This viewpoint is in fact not surprising since in the color evaporation model \[17\] based on duality arguments, the cross section for charmonium production is written as an integral over the $s$ range $4m_c^2$ and $4m_{D^0}^2$ (i.e. the upper limit equal to open charm production). However, charmonium production takes place in a “hot” medium as Hoyer pointed out in his talk (for more details see Ref. \[18\]), so sub-thresholds effects likely should have an influence. In the improved color evaporation model this means that the heavy quark pair can get energy from the surrounding color field to give open charm production beneath $s = 4m_{D^0}^2$. \[19\]

Analogously in the COM, those dipole transitions occurring during the evolution into final particles should be associated to absorption of soft gluons from the neighborhood as well. For example a $c\bar{c}(3S_1^{(8)})$ pair may absorb two non-perturbative gluons becoming a $O(v^2)$ Fock component of the $J/\psi$ state or may radiate two soft gluons thus with nearly unity overlap with the $J/\psi$ wave function. The conclusion is that on the average one should consider both possibilities, likely decreasing the net impact of the above-mentioned effect on the $p_t$ distribution.

4. LHC PREDICTIONS

As outlined in the previous section, the extraction of NRQCD matrix elements from charmonium production is not yet fully satisfactory due to several theoretical uncertainties. Still more analysis has to be done to check the universality hypothesis allowing comparisons between different production reactions. Polarization of charmonium from gluon fragmentation at high $p_t$ might be decisive to verify the validity of the COM \[20\].

Nevertheless, under the assumption of the validity of the COM to explain charmonia production, long-distance matrix elements obtained from Tevatron data should parameterize those uncertainties to a large extent. Therefore, in extrapolating to LHC energies, predictions should be reliable if correctly performed. In our case, this actually means running PYTHIA with the same parameters and options (e.g. radiation on) as previously employed in the fit to Tevatron experimental points.

In Figure 2 we show our predictions for prompt $J/\psi$ direct production at the LHC using the CTEQ PDF. A caveat is in order, however. In changing the choice for the PDF there are noticeable changes on the production rate by about 50% in certain regions of $p_t$ (see Ref. \[17\]).

4.1. Muon cross sections from prompt $J/\psi$ production at the LHC

We have estimated the cross sections for events with muons from prompt $J/\psi$’s, passing typical pseudorapidity and transverse momentum cuts, foreseen in the first and second level triggers for $B$ physics at the two LHC experiments ATLAS and
Figure 2. Our predictions for prompt $J/\psi$ direct production at the LHC according to the color-octet model implemented in PYTHIA for the CTEQ PDF [10]. Dotted line: CSM; dashed line: $^{1}S_{0}^{(8)} + ^{3}P_{J}^{(8)}$ contributions; dot-dash line: $^{3}S_{1}^{(8)}$ contribution; solid line: all contributions. Notice that asymptotically the $^{3}S_{1}^{(8)}$ contribution becomes dominant, likely implying a transverse polarization of charmonium at high $p_t$.

CMS. We shall denote as $\mu 3$, $\mu 6$ ($\mu 3\mu 3$, $\mu 6\mu 3$) those events with at least one (two) muon(s) whose $p_t$ is (are) larger than 3 GeV or 6 GeV, respectively.

From our modified version of PYTHIA with the COM implemented in, we obtain

$$\sigma(\mu 3) = 622 \text{ nb}$$
$$\sigma(\mu 6) = 52 \text{ nb}$$
$$\sigma(\mu 3\mu 3) = 30 \text{ nb}$$
$$\sigma(\mu 6\mu 3) = 9 \text{ nb}$$

In all cases an additional pseudorapidity cut $|\eta| < 2.5$ for muons was required on generated events.

Those numbers should be useful to evaluate an order-of-magnitude estimate of background sources for the $B$ physics programme at LHC experiments, especially regarding the search for CP violation via the golden decay mode $B_{d}^{0} \rightarrow J/\psi K_{s}$.

Acknowledgments

I thank P. Eerola, N. Ellis, P. Hoyer and M. Smizanska for interesting discussions.

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