Attachment of Tripping Wires to Enhance the Efficiency of a Vortex-Induced Vibrations Energy Generation System*

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Abstract
We conducted an experiment to quantify the increase in efficiency of energy extraction from a constant flow that can be obtained by attaching a pair of tripping wires to an elastically mounted cylinder under influence of vortex-induced vibrations. Free oscillation tests were carried with three different cylinder configurations: smooth, with tripping wires positioned at angular positions equal to 60° and 75°. The Reynolds number varied from 2.9×10³ to 2.2×10⁴. We measured the amplitude of oscillation and the output voltage to calculate the power generated and conversion efficiency. The maximum power generation occurred when tripping wires were positioned at 60° and reduced velocity was 12, but only a 2.88% efficiency was achieved at this case. The maximum efficiency obtained was 12.47% and occurred when tripping wires were positioned at 75°, with reduced velocity 6.5. The maximum efficiency with tripping wires attached was about four times larger than the maximum efficiency obtained on smooth cylinder case. The presence of tripping wires also widen the lock-in region and when they were placed at 75° an almost constant power generation was obtained within reduced velocity ranging from 6 to 8.5, meaning that it is possible to design a system that generates the same amount of power even with variations on flow speed.

Key words: Vortex-Induced Vibrations, Renewable Energy, Tripping Wires, Energy Harvesting, Circular Cylinder

Nomenclature

| Symbol | Definition | Value and Unit |
|--------|------------|----------------|
| A      | Non-dimensional amplitude of oscillation | [-] |
| A_peak | Displacement signal peak amplitudes | [m] |
| d      | Tripping wires diameter | [m] |
| D      | Cylinder diameter | [m] |
| E      | Generated Voltage | [Volt] |
| f_n   | System natural frequency of oscillation | [Hz] |
| f_osc | Frequency of oscillation | [Hz] |
| f_s   | Frequency of vortex-shedding | [Hz] |
| L      | Cylinder length | [m] |
| m      | System oscillating mass | [kg] |
| m*     | Mass ratio | \[\frac{m}{\rho \pi d^2 L}\] |
| P      | Electric power generated | [kg m² s⁻³] |
| Re     | Reynolds Number | [-] |
1. Introduction

The world’s energy supply is hardly dependent of fossil fuels. The World Energy Outlook 2007 predicted that coal, oil and gas will be responsible to meet 84% of energy demand in 2030 and it is of common knowledge that these non-renewable sources will deplete in a near future, however it is not a consensus when it will happen because lots of efforts have been made to delay it. The search for new oil fields is constant, as well the development of new drilling and extraction techniques to increase the volume of recovered oil. In parallel, the energy conversion efficiency of existing systems is being improved and renewable energy sources are being developed. The latest World Energy Outlook, released in 2012, estimates that is necessary to invest 4.8 trillion US dollars on the development of renewable energy systems until 2035 to meet the demand increase.

A promising source of renewable energy is the ocean. There are two basic principles on producing energy from it that are usually broached. The first one is thermal and relies on the difference of temperature between the surface and deep water. The warm water on the surface is used to vaporize a fluid, which has a low boiling point, and the vapor expansion turns a turbine that is connected to a generator, producing electricity. Then the vapor is directed deep into the ocean where it condenses, re-starting the cycle. The second one consists on using the mechanical energy of tides and waves. It is based on the principle that a floating body moves up and down under effect of tides and waves. The body is connected to a linear type generator that converts its vertical motion in electricity, but this kind of generation can be intermittent as depends of weather conditions.

There is a third potential source of energy on the oceans that is not explored: the currents. The main reason is because ocean currents have low speeds, usually less than 1.0 m/s at water surface, and it is not efficient to use turbine type generators at such low speed. Recently, a group of researchers of Michigan University developed a system called VIVACE\textsuperscript{1} (Vortex-Induced Vibration Aquatic Clean Energy) that consists of an elastically mounted cylinder that is submerged on water and connected to a generator. Their system takes advantage of a phenomenon called vortex-induced vibrations to extract energy from a flow.

Vortex-induced vibrations are the result of vortex shedding that occurs on a bluff body under influence of a steady flow. The vortex shedding process induces periodic forces on the body which can result in motions. The motions are perpendicular to the flow direction and are self-excited when the vortex-shedding frequency and system natural frequency are approximately the same. The vortex shedding frequency can be predicted using the Strouhal number equation:
The strouhal number of a circular cylinder is well known and is almost constant around 0.21, at Reynolds numbers greater than $1\times10^3$. Once known the vortex shedding frequency it is only necessary to set the system natural frequency equal to the vortex shedding frequency to obtain motions on the cylinder. There are other factors that influence the amplitudes of oscillation, as system damping and mass ratios, and Reynolds number, but further discussion is out of the current work scope.

Following VIVACE\(^{1}\) steps, Nishi et al.\(^{4}\) developed a system that utilizes the vortex-induced vibrations of a cylinder and principle of leverage. Their system transforms the translation motion of the VIV into a translation motion of the generator as shown in Fig. 1. The aim of their work was to improve the system efficiency by tuning the distance between the pivot and the generator, changing the relation between the cylinder motion and coil motion.

The efficiency of energy conversion using vortex induced vibration systems needs to be improved. VIVACE\(^{1}\), the most efficient system reported until now, achieves only 22%. The search of a technique to improve the conversion efficiency led to study the work of researchers\(^{(5)-(10)}\) who have studied the addition of attachments to the cylinder in order to change the flow characteristics.

Igarashi\(^{(5)}\) and Nebres and Batil\(^{(6)}\) concentrated efforts on measuring hydrodynamic forces and strouhal number of cylinders with tripping wires attached at angular positions $\beta$, defined in Fig. 2. They found three flow patterns, shown in Fig. 3, which are dependent of Reynolds Number and tripping wire diameter and angular position. The first pattern (A), shown in Fig. 3.a, is characterized by a separation over tripping wire, an immediately reattachment behind it and a posterior laminar re-separation, not causing changes on drag coefficients. The second pattern (B), shown in Fig. 3.b, has as main characteristic a separation over the tripping wire, a reattachment behind it and a turbulent re-separation that occurs more downstream than the smooth cylinder case, resulting in a narrow wake region and, consequently, lower drag coefficients than smooth cylinder. The third and last pattern (C), shown in Fig. 3.c, is characterized by a completely separation over the tripping wire, resulting in a wider wake region. The widening on the wake region led to the formation of larger vortexes behind the cylinder, which in turn led to larger pressure fluctuations and, consequently, larger drag and lift forces.

\[ St = \frac{f_c D}{V} \]  

(1)
Alam et al.\textsuperscript{(7)} and Quadrante and Nishi\textsuperscript{(8)} studied the effects of a pair of tripping wires positioned symmetrically in relation to the stagnation point. Alam et al.\textsuperscript{(7)} focused the reduction of drag and lift forces acting on a non-oscillating cylinder with tripping wires positioned within $20^\circ < \beta < 60^\circ$. The comparison between tripping wires and smooth cylinder cases showed two distinct regions. They found that when $20^\circ < \beta < 40^\circ$ the forces acting on the cylinder decreased but, when $50^\circ < \beta < 60^\circ$, the forces increased. They also confirmed the flow patterns B and C described previously as the cause of the decrease and increase of forces, respectively. Quadrante and Nishi\textsuperscript{(8)} expanded tripping wires positioning range to $15^\circ < \beta < 90^\circ$ to obtain more data about the increase on forces and included free oscillating tests to obtain VIV response with tripping wires. They found large increase on forces and VIV amplitudes when $\beta = 60^\circ$ and $\beta = 75^\circ$. As the experimental apparatus of their work and the present work are basically the same their results are of great relevance, then a revised data analysis of their experiments is presented briefly in a subsequent section.

![Flow patterns related to tripping wires position.](image)

The present study has the objective of combining Nishi et al.\textsuperscript{(4)} and Quadrante and Nishi\textsuperscript{(8)} studies to achieve a more efficient system by exploring the effects present on pattern C which causes larger fluid forces on the cylinder. To induce the occurrence of pattern C a pair of tripping wires was attached to the cylinder at $60^\circ$ and $75^\circ$ and free oscillation tests were carried out to measure the amplitude of oscillation, the forces and the voltage generated. The results were compared with smooth cylinder case to quantify tripping wires contribution to system efficiency. The next two sections contain a brief description of the mathematical model behind the principle of leverage and a revision of tripping wires effects on VIV response obtained by Quadrante and Nishi\textsuperscript{(8)}. They are followed by the description of current experimental setup, the presentation of results and the


2. Mathematical model

Table one contains the variables and parameters pertinent to the mathematical model. The leverage system is represented by a single rigid bar that rotates around a pivot, as shown on Fig. 4. Fluid force from VIV and electromagnetic reaction forces are applied at distances $r_2$ and $r_3$, respectively, from the pivot, while spring restoring force is applied at distance $r_1$. The dynamic equation of the system is written as:

$$I\ddot{\theta} + C\dot{\theta} + k\theta + kr_1^2 \sin\theta \cos\theta = F_Lr_2 \cos\theta + Mr_3 \cos\theta - F_mr_3 \cos\theta$$

(2)

Fig. 4: System modeled as a single rigid bar.

The voltage generated on the generator is proportional to the velocity of the coil’s translation motion and can be written as:

$$E = \alpha r_1 \dot{\theta} \cos\theta$$

(3)

Where $\alpha$ is the constant of proportionality obtained experimentally. The electromagnetic reaction force is:

$$F_m = \frac{am_br_1N}{\ell R} \theta \cos\theta$$

(4)

By solving the equation of motion and considering the frequency of lift force equal to the system’s natural frequency it is possible to obtain the equation of maximum generated voltage:

$$E_{max} = \frac{\alpha F_L}{C + \frac{am_br_1N}{\ell R} \frac{r_2}{r_3}}$$

(5)

Table 1: Parameters and variables pertinent to numerical model

| Symbol | Definition | Value |
|--------|-----------|-------|
| C | Damping coefficient | $2\zeta \sqrt{k}$ [kg m s$^{-1}$] |
| $F_L$ | Amplitude of lift force | - [N] |
| $F_m$ | Electromagnet reaction force | - [N] |
| $g$ | Gravity acceleration | 9.81 [kg m s$^{-2}$] |
| $I$ | Moment of Inertia of the bar and coil | $0.48 + Mr_r^2$ [kg m$^2$] |
| $k$ | Spring stiffness | 540 [N m$^{-1}$] |
| $\ell$ | Coil length | 0.043 [m] |
| $m_b$ | Magnet bar coefficient | $0.637 \times 10^{-3}$ [Wb] |
| $M$ | Coil mass | 0.548 [kg] |
| $N$ | Total number of turns of the coil | 2000 |
| $R$ | Resistor connected to the generator | 100 [Ω] |
### 3. Tripping wires effects

A revised data analysis of Quadrante and Nishi\(^{(8)}\) experiments with tripping wires is presented on this section. They conducted non-oscillating tests with a low mass and damping ratios system \((m^*=6.94\text{ and } \zeta=0.016, \text{ respectively})\) to determine the position of maximum increase on lift coefficient and chose two cases to use on free oscillating tests. Figure 5 shows lift and drag coefficients obtained from non-oscillating tests in function of tripping wires angular position. The maximum lift and drag coefficients occur when \(\beta = 75^\circ\) and are, respectively, 65% and 44% larger than smooth cylinder case.

![Figure 5: Lift and drag coefficients of a non-oscillating cylinder in function of tripping wires position (Quadrante and Nishi\(^{(8)}\))](image)

4. Experiment

#### 4.1 Experimental arrangement

The experiments were conducted in a 100m long, 8m wide and 3.5m depth towing tank at Yokohama National University. The system arrangement, presented in Fig. 7, was built of aluminum bars connected by bearings. Its construction guarantees an almost perpendicular movement of the cylinder as the arm length measures 0.5m and the maximum vertical oscillation expected is 0.025m. The cylinder motion can be approximated as a pure translation motion in vertical direction as the maximum horizontal excursion expected is \(6.25\times10^{-4}\text{ m}\), much smaller than the vertical.
The cylinder also was built of aluminum with external diameter $D = 0.025m$ and length $L = 0.454m$. It was supported by four springs in parallel, resulting in stiffness $k = 540N/m$. Two bi-axial load cells were equipped at cylinder ends to measure the drag and lift forces acting on it. To measure the vertical motion, a laser displacement sensor was positioned right over the cylinder. Acrylic endplates were attached to cylinder extremities in order to avoid end effects and ensure a two-dimensional flow.

The tripping wires in this study were the same used by Quadrante and Nishi and a schematic view of their positioning around the cylinder is shown in Fig. 8. They are placed symmetrically in relation to the stagnation point at angular positions $\beta$. The diameter $d = 0.003m$ results in a diameter ratio $d/D = 0.12$. The gap $\delta$ between tripping wires and cylinder surfaces were equal to 0.0025, resulting in a gap ratio $\delta/D = 0.1$.

The electricity generator attached to the system is composed of a fixed magnetic bar and a moving coil. A 100Ω resistor is connected to the generator to act as a load and the voltage on it is measured.
4.2 Methodology

Free oscillation tests were carried out with and without tripping wires to compare amplitude response and generated voltage with nine different configurations which are combinations of three cylinder characteristics (smooth cylinder, tripping wires at $\beta = 60^\circ$ and $\beta = 75^\circ$) and three generator positions ($r_3/r_2 = 0.9, 1.0$ and $1.1$). The relation $r_3/r_2$ tested was determined by substituting the system parameters in Eq. 5 and varying $r_3$ to create a maximum voltage response in function of $r_3$. The result is shown in Fig. 9, from which is possible to note that the maximum voltage is obtained when $r_3 = r_2$.

![Tripping wires positioning](image)

Fig 8: Tripping wires positioning

![Theoretical maximum voltage response in function of $r_3/r_2$.](image)

Fig 9: Theoretical maximum voltage response in function of $r_3/r_2$.

The damping ratio ($\zeta$) of the system is composed by the structural damping ($\zeta_{\text{structural}} = 0.010$), measured without the magnet positioned inside the coil, and the generator damping ($0.038 \leq \zeta_{\text{generator}} \leq 0.042$), calculated by subtracting the structural damping from the total damping ratio, which was obtained from free-decay tests with the $100\Omega$ resistor attached to the generator. Figure 10 contains one example of the free-decay tests signal with the exponential curve (Eq. 6) adjusted to the selected peaks from the signal. The value of $\zeta$ used to the best adjustment of the equation is the damping ratio of the system. The linear approximation shows a good agreement with the signal inside the selected region, in red. The selected region of the signal was made by considering the maximum amplitudes of oscillation expected of 0.025m.

The decay tests were also used to obtain the system natural frequency of oscillation in water. The natural frequency is essential to calculate the towing speed in order to cover the desired reduced velocity range, $2 < V_R < 12$, within vortex induced vibrations of circular cylinder are expected. Inside this reduced velocity range, Reynolds number varied by $2.9 \times 10^3 < Re < 2.2 \times 10^4$. At least two repetitions of each reduced velocity for each case were made to check the repeatability of VIV response.
\[ Z = Z_0 e^{-\zeta^2 \omega_0 t} \]  \hfill (6)

Fig 10: Free-decay test response with the exponential curve adjusted to the peaks

4.3 Summary of experiment parameters

All the system parameters that are related to the flow are summarized in table 2.

| Parameter                      | Value and Unit |
|--------------------------------|----------------|
| Cylinder diameter \((D)\)       | 0.025 [m]      |
| Cylinder Length \((L)\)         | 0.454 [m]      |
| Tripping wires diameter \((d)\)| 0.003 [m]      |
| Range of Reynolds number \((Re)\)| \(2.9 \times 10^3 < Re < 2.2 \times 10^4\) [-] |
| Range of reduced velocity      | \(2 < V_R < 12\) [-] |

4.4 Data processing

We measured the drag and lift forces acting on each cylinder end, the towing carriage speed, the cylinder vertical displacement and the voltage at the resistor. The first step of data analysis was the zero adjustment and digital filtering of all signals, followed by the selection of the region for temporal and spectral analysis.

Total forces were obtained by the sum of forces measured by load cells. Drag coefficient was calculated using the average drag force \((\bar{F}_D)\) and the lift coefficient was calculated using the fluctuating lift force \((F_{Lrms})\) as follows:

\[ C_D = \frac{2\bar{F}_D}{\rho DLV^2} \]  \hfill (7)

\[ C_L = \frac{2F_{Lrms}}{\rho DLV^2} \]  \hfill (8)

Non-dimensional amplitudes were calculated from the maximum absolute amplitudes \((A_{\text{peaks}})\) of cylinder displacement. A peak detection routine written in Matlab® was used to find all signal peaks and then compute the average of the peaks. The amplitude was normalized by the cylinder diameter as usual in VIV analysis as:

\[ A = \frac{1}{n} \sum_{i=1}^{n} A_{\text{peaks}} \]  \hfill (9)

All signals dominant frequencies and respective phases were obtained through Fast Fourier Transform (FFT). The relative phase between the displacement and lift force \((\phi_F)\) is calculated as:

\[ \phi_F = \phi_L - \phi_Z \]  \hfill (10)
Where $\phi_Z$ and $\phi_L$ denote the displacement and lift force phases, respectively. This definition of $\phi_L$ means that the lift force excites the vibration when $\phi_L$ ranges from 0 to 180°, while the lift force decays the vibration when $\phi_L$ ranges from -180° to 0°.

Energy conversion efficiency is calculated as the ratio between electrical power $P$ and fluid flow power $W_{\text{water}}$:

$$\eta = \frac{P}{W_{\text{water}}} = \frac{1}{0.5\rho V^3 DL} \int_{T_f - T_s}^{T_f} \frac{E^2}{R} dt$$

(11)

5. Results

The three variations of $r_3/r_2$ presented similar results, with small variations in maximum amplitudes of VIV, voltage and generated power. As expected from Fig. 9, the maximum generated voltage was obtained for $r_3/r_2 = 1.0$ and for that reason we chose this case to present in this section. A summary containing all the results, including $r_3/r_2 = 0.9$ and 1.1, will be presented on the end of this section.

5.1 VIV amplitude response

Figure 11 contains the VIV amplitude and frequency responses. Both cases with tripping wires present larger amplitudes and a wider lock-in region than smooth cylinder case. The maximum normalized amplitudes for $\beta = 75^\circ$ and $\beta = 60^\circ$ are, respectively, 0.63 and 0.74, however, for smooth cylinder case, it is only 0.23. The monotonically increase in amplitude related by Quadrante and Nishi(8) and presented in Fig. 6 can be observed on Fig. 11, but it occurred only when $\beta = 60^\circ$ and $V_R \geq 10$ and, even then, not for all repetitions. For $V_R = 10$ and 11 it is possible to see distinct responses, with one of the test runs presenting large amplitudes and the other one a near zero value.

The monotonically increase of amplitude present in Fig. 6 is related to a transition from VIV to galloping phenomenon, which occurs at higher reduced velocities than VIV. The galloping phenomenon is the result of wake instabilities which induce lift forces on the cylinder. The difference in the galloping response between Figs. 6 and 11 are associated to two factors: difference in the mass and damping ratios between experiments and the position of tripping wires.

Firstly the difference in mass and damping ratios will be discussed. The current experiment was carried with larger mass ratios (10.08 ≤ $m^*$ ≤ 11.28) and damping ratios (0.042 ≤ $\zeta$ ≤ 0.052) than the presented by Quadrante and Nishi(8) ($m^*$=6.94 and $\zeta$=0.016). The lower values in their experiments facilitate the occurrence of large galloping.

The second factor that contributes for galloping occurrence is the tripping wires position. Chang et al.\textsuperscript{(10)} studied the effects of surface roughness in VIVACE\textsuperscript{(1)} galloping response and found that systems with large mass and damping ratios are susceptible to galloping when the roughness is placed inside a critical range of angular positions, from 10° to 62°. Above 62°, the placement of roughness inhibited the galloping phenomena. This explains the non-existent galloping when tripping wires were placed at $\beta=75^\circ$. The unstable occurrence of galloping when tripping wires were at $\beta=60^\circ$ is result of the proximity to 62° limit observed by Chang et al.\textsuperscript{(10)}, which is a transition region.

The lock-in range, characterized by $f_{\omega_{\text{osc}}}/f_n \approx 1$, also presents differences between the three cases. Smooth cylinder and $\beta = 75^\circ$ cases have a limited lock-in range while $\beta = 60^\circ$ lock-in range starts at $V_R = 6.0$ and extends until the maximum reduced velocity tested.
To understand the changes caused on amplitude and lock-in we analyzed the lift coefficients and the difference of phase between lift force and cylinder displacement (Eq. 10) shown in Fig. 12. To achieve large amplitude response a favorable combination of lift force and phase difference is necessary. According to the definition in Eq. 10, a phase difference equal to 90° means that the lift force acts in a purely excitation rule. It means that as near the difference of phase is to 90° more efficiently the force excites the system and larger amplitudes of oscillation can be expected.

Smooth cylinder case presents lower lift coefficients than tripping wires cases, except at the region between reduced velocities 4.0 and 5.5 within it achieved almost the same maximum coefficient than \( \beta = 75^\circ \) case. However, even these large lift coefficients were not translated into high amplitudes because the phase difference is around 30°. If we compare the smooth cylinder to the lock-in region \( 6 \leq V_R \leq 8.5 \) of the case with tripping wires at 75° it is noticeable the rule of the phase difference on VIV response. This case has similar maximum lift coefficients than the smooth cylinder, but the oscillation amplitudes achieved were more than three times larger than the smooth cylinder case, just because the difference of phase is around 70°.

The same principle can be used to explain the lower amplitudes achieved when tripping wires were positioned at \( \beta = 60^\circ \) than at \( \beta = 75^\circ \). The lift coefficients for \( \beta = 60^\circ \) are slightly lower than the smooth cylinder, as well the phase, around 60°. The composition of smaller coefficients with a phase difference more distant from 90° led to smaller amplitude. The opposite occurred at reduced velocities larger than 9, when the case of \( \beta = 60^\circ \) presented much larger amplitudes than the \( \beta = 75^\circ \) case. Now, the lift coefficients are a little larger for \( \beta = 60^\circ \) and the phase difference stills around 60°, while the phase for \( \beta = 75^\circ \) shifted to near 180°.

5.2 Generated power and conversion efficiency

Generated power was calculated using the numerator of Eq. 11 and is shown in Fig. 13. The maximum electric power generated was 57.8mW and occurred on \( \beta = 60^\circ \) case at \( V_R = 12 \). \( \beta = 75^\circ \) case also presented large electric power between reduced velocities 6.5 and 8.5, achieving up to 50.41mW. Smooth cylinder case had the lowest power generation, achieving a maximum of 4.86mW at \( V_R = 6.0 \).
Fig. 12: Lift coefficients and difference of phase between lift force and cylinder vertical position. Symbols: (○) smooth cylinder, (□) tripping wires at $\beta = 60^\circ$, (▲) tripping wires at $\beta = 75^\circ$.

Figure 14 shows the efficiency in function of reduced velocity. $\beta = 75^\circ$ case had the best performance, achieving a maximum efficiency $\eta = 12.41\%$ at $V_R = 6.5$, about 4 times the maximum efficiency obtained on smooth cylinder case. Although the maximum power generation occurred in $\beta = 60^\circ$ case its efficiency is very low because the high flow speed involved.

Fig. 13: Electric power generated. Symbols: (○) smooth cylinder, (□) tripping wires at $\beta = 60^\circ$, (▲) tripping wires at $\beta = 75^\circ$.

Fig. 14: System conversion efficiency. Symbols: (○) smooth cylinder, (□) tripping wires at $\beta = 60^\circ$, (▲) tripping wires at $\beta = 75^\circ$. 
5.3 Results summary

Table 3 contains a summary of all results obtained. The maximum amplitude response and electric power generated occur when \( \frac{r_3}{r_2} = 1.0 \) and \( \beta = 60^\circ \), at \( V_R = 12 \), but this case presents a low efficiency because the power of the flow is high at this speed. Smooth cylinder cases presented very small amplitude response and, consequently, low electric power generation and efficiency at all tested configurations. The most efficient case was \( \beta = 75^\circ \), presenting values between 10.71% and 12.47%, at reduced velocities between 6 and 8.

From Table 3 is also possible to realize that the presence of large lift forces acting on the cylinder do not guarantee a high amplitude response. The phase difference between lift force and displacement also has an important influence since as near it is to \( 90^\circ \), larger will be the force contribution to excite the system VIV.

6. Discussion

The efficiency of energy generation systems using vortex-induced vibrations has been studied in a few works\(^{(1)(4)}\). The main characteristics and results of previous and current works are listed in Table 4. VIVACE\(^{(1)}\) is the best system developed until the present moment, achieving a 22% conversion efficiency, value much larger than obtained by Nishi et al.\(^{(4)}\) and by the current experiment with smooth cylinder.

| Test configuration | Natural frequency (Hz) | Mass and Damping ratios | Maximum Value | Max. Lift Coefficient | Max. Electric Power (mW) | Max. Efficiency (%) |
|--------------------|------------------------|-------------------------|----------------|-----------------------|-------------------------|-------------------|
| \( r_3/r_2 \) case | \( f_n \) | \( m^* \) | \( \zeta \) | \( V_R \) | \( q_y(\circ) \) | \( V_R \) | \( V_R \) | \( V_R \) |
| Smooth             | 2.44                   | 9.88                    | 0.041          | 6.0                   | 0.35                     | 41                | 0.84             | 5.5               | 15.21            | 6.0               | 2.48             | 6.0               |
| \( \beta = 60^\circ \) | 2.42                   | 10.08                   | 0.050          | 34                    | 0.44                     | 34                | 0.55             | 7                  | 26.19            | 7.0               | 6.08             | 7.0               |
| \( \beta = 75^\circ \) | 2.42                   | 10.08                   | 0.045          | 72                    | 0.61                     | 72                | 0.96             | 6                  | 55.07            | 8.0               | 10.71            | 7.0               |
| Smooth             | 2.38                   | 10.38                   | 0.042          | 35                    | 0.23                     | 35                | 0.92             | 5.5               | 4.87             | 6.0               | 2.82             | 6.0               |
| \( \beta = 60^\circ \) | 2.36                   | 10.60                   | 0.048          | 59                    | 0.74                     | 59                | 0.58             | 12                 | 57.70            | 12.0              | 5.12             | 6.5               |
| \( \beta = 75^\circ \) | 2.36                   | 10.60                   | 0.045          | 74                    | 0.63                     | 74                | 1.03             | 6                  | 50.47            | 8.0               | 12.41            | 6.5               |
| Smooth             | 2.30                   | 10.10                   | 0.058          | 13                    | 0.15                     | 13                | 0.75             | 5                  | 0.45             | 5.5               | 2.21             | 5.5               |
| \( \beta = 60^\circ \) | 2.29                   | 11.13                   | 0.053          | 31                    | 0.67                     | 31                | 0.59             | 12                 | 50.28            | 12.0              | 5.96             | 6.5               |
| \( \beta = 75^\circ \) | 2.28                   | 11.28                   | 0.052          | 69                    | 0.55                     | 69                | 0.86             | 6.5               | 46.99            | 7.5               | 12.47            | 6.5               |

There are two main differences between the VIVACE\(^{(1)}\) and the other two systems: the mass ratio and Reynolds number. Low mass ratio can increase the length of lock-in region, as discussed by Khalak and Williamson\(^{(2)}\), but has small influence on amplitude response. On the other hand, variation of Reynolds number has large influence on response because the regime of vortex generation changes\(^{(3)}\) as it increases. Nishi et al.\(^{(4)}\) and the current experiment operates in lower Reynolds number, under TrSL2\(^{(11)}\) (transition of shear layer 2) regime, which is a subcritical regime characterized by the formation of transition vortices in free shear layer. VIVACE\(^{(1)}\) operates on TrSL3\(^{(11)}\) regime, which also is a subcritical regime, but has a fully turbulent shear layer that results in larger fluid forces and amplitude response. The attachment of tripping wires to our system increased the lock-in range and amplitudes of oscillation without necessity of change on Reynolds number. It probably induced a transition from TrSL2 to TrSL3 regime even at lower Reynolds number, but it is not possible to confirm without obtaining the vortex pattern behind the cylinder.

The system efficiency had a significant increase on efficiency when \( \beta = 75^\circ \). Moreover, the amplitude of oscillation and, consequently, the power generated had almost stabilized values within the lock-in range \( (6.0 \leq V_R \leq 8.5) \). This characteristic is very important as it is
expected that the current speed suffer fluctuations along time and, with a large lock-in range, it is possible to tune the system to deal with these fluctuations to minimize downtime.

Table 4: Characteristics and efficiency of known energy generation systems that uses VIV.

| Experiment          | Re    | $m^*$ | $\zeta$ | $\eta$ (%) |
|---------------------|-------|-------|---------|------------|
| VIVACE              | 9.4x10^4 | 1.45  | 0.059   | 22.00      |
| Nishi et. al.       | 5~7x10^3  | 12.68 | 0.023   | 4.50       |
| Current experiment  |       |       |         |            |
| Smooth $\beta = 60^\circ$ | 2.9~22x10^3 | 10.08 | 0.050   | 6.08       |
| $\beta = 75^\circ$  | 11.28 | 0.052 | 12.47   |            |

By observing Table 4 it is clear that lots of new configurations can be studied to search for a more efficient system. It is not possible to predict what will happen if tripping wires are attached to VIVACE, or what will happen if the current experiment is conducted at higher Reynolds numbers. Another points are that the generator of the current experiment is not optimized to work with the force increase resulted from the attachment of the tripping wires and there was no investigation of the effects of changing the load required form the generator in the system response. Changes in the load causes significant changes in the damping of the system and must be investigated.

Further research and experiments are necessary to clarify these questions and check the real potential of generating energy from a fluid flow using vortex-induced vibrations phenomenon.

7. Conclusion

This study carried out experiments about the attachment of a pair of tripping wires to the cylinder of an energy extraction system through vortex-induced vibrations and use principle of leverage to transmit the oscillations from the cylinder to the generator. We attached tripping wires at two different angular positions, $\beta = 60^\circ$ and $\beta = 75^\circ$, and conducted free oscillating experiments to measure response amplitude and power generated to compare with smooth cylinder results.

We found that the presence of tripping wires widen the lock-in range and increased the amplitudes of oscillation and power generated significantly. At the best configuration, when $\beta = 75^\circ$, the system efficiency achieved $\eta = 12.47\%$, more than four times the efficiency obtained on smooth cylinder experiments, showing that the use of tripping wires has a great potential to be used to enhancement of vortex-induced vibrations.

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