One-dimensional Repulsive Fermi Gas in a Tunable Periodic Potential

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By using unbiased continuous-space quantum Monte Carlo simulations, we investigate the ground state properties of a one-dimensional repulsive Fermi gas subjected to a commensurate periodic optical lattice (OL) of arbitrary intensity. The equation of state and the magnetic structure factor are determined as a function of the interaction strength and of the OL intensity. In the weak OL limit, Yang’s theory for the energy of a homogeneous Fermi gas is recovered. In the opposite limit (deep OL), we analyze the convergence to the Lieb-Wu theory for the Hubbard model, comparing two approaches to map the continuous-space to the discrete-lattice model: the first is based on (non-interacting) Wannier functions, the second effectively takes into account strong-interaction effects within a parabolic approximation of the OL wells. We find that strong antiferromagnetic correlations emerge in deep OLs, and also in very shallow OLs if the interaction strength approaches the Tonks-Girardeau limit. In deep OLs we find quantitative agreement with density matrix renormalization group calculations for the Hubbard model. The spatial decay of the antiferromagnetic correlations is consistent with quasi long-range order even in shallow OLs, in agreement with previous theories for the half-filled Hubbard model.

Making unbiased predictions for the properties of strongly correlated Fermi systems is one of the major challenges in quantum physics research. One-dimensional systems play a central role in this context since, on the one hand, correlations effects are more pronounced in low dimensions and, on the other hand, exact results have been derived in a few relevant cases. Two such cases are the homogeneous Fermi gas, whose exact ground-state energy was first determined by Yang via the Bethe Ansatzt technique, and the single-band Hubbard model, whose solution was provided by Lieb and Wu. These two paradigmatic models describe two opposite limits of realistic physical systems, which in general are neither perfectly homogeneous nor devoid of interband couplings. In the absence of exact analytical theories for the more realistic intermediate regime, developing unbiased computational techniques is of utmost importance.

The experiments performed with ultracold atoms trapped in optical lattices (OLs) have emerged as the ideal playground to investigate quantum many-body phenomena in periodic potentials. The intensity of the external periodic field can be easily varied by tuning a laser power, and also the interaction strength can be tuned exploiting Feshbach resonances. This has recently allowed the remarkable observation of antiferromagnetic correlations in a controlled experimental setup, both in two and in one dimension. The bulk of early research activity on OL systems focussed on deep OLs and weak interactions, where single-band tight-binding models are adequate. Away from this regime multi-band processes come into play, and the effect of the independent tuning of the OL intensity and the interaction strength can be captured only via multi-band or continuous-space models. Recent theoretical and experimental studies have addressed the regime of shallow OLs and strong interactions, investigating intriguing phenomena such as Mott and pinning bosonic localization transitions, Anderson localization, Bose-Glass phases, and itinerant ferromagnetism.

Previous theoretical studies on extended one-dimensional Fermi gases considered either homogeneous continuous-space systems or discrete-lattice models. In this Rapid Communication, we investi gate the ground-state properties of a continuous-space one-dimensional Fermi gas with zero-range repulsive interactions, subjected to a periodic potential (representing an OL) of arbitrary intensity. We focus on a balanced (i.e., unpolarized) two-component mixture at the density of one fermion per well (half filling). The energy and the magnetic structure factor are computed via continuous-space diffusion Monte Carlo (DMC) simulations, which provide unbiased predictions for one-dimensional Fermi systems.

We explore the crossover between two opposite limits. For a vanishing OL, we recover the ground state energy of a homogeneous system predicted by Yang; for a deep OL, where the continuous-space system can be mapped to a discrete-lattice model, we inspect the convergence to the Lieb-Wu results for the Hubbard model. Specifically, we consider two mapping procedures; the first is based on the standard Wannier functions, the second is designed to effectively take into account within an harmonic approximation the higher-orbital effects induced by strong interactions. The regimes where these two mapping procedures become quantitatively accurate are outlined. Furthermore, the onset of the antiferromagnetic correlations is explored.
The external potential spin-down particles). The index particles (irrespectively of their spin state), while the polarized) mixture of the two components only spin-down particles. We focus on a balanced (unpolarized) ground state interaction energy per particle \( E_{\text{int}}/N = (E - E_{\text{gamma}})/N \) as a function of the interaction parameter \( \gamma = 2/(n|a_{1D}|) \). \( E \) and \( E_{\text{gamma}} \) are the energies of an interacting and a noninteracting gas in an OL, respectively. The density is fixed at half filling \( n = N/(Ld) = 1/d \). Symbols connected by dashed lines correspond to QMC results (system size \( L = 26 \)) for different OL intensities \( V \), expressed in units of the recoil energy \( E_r \). The thick continuous curve is the Yang’s Bethe-Anstatz result [2] for the homogeneous Fermi gas (\( V = 0 \)).

that strong correlations form in deep OLs, where the continuous-space DMC data agree with Hubbard-model results, which we obtain using the density matrix renormalization group (DMRG) method. Interestingly, we find that the correlation amplitude can be large even in very shallow lattices if the interaction strength is tuned close to the infinite repulsive (Tonks-Girardeau) limit. Both in deep and in shallow OLs the spatial decay of the correlations appears to be consistent with the quasi-long range order predicted by bosonization theories for the half-filled Hubbard model and for the one-dimensional Wigner crystal.

We consider a one-dimensional two-component atomic Fermi gas described by the following continuous-space Hamiltonian:

\[
\hat{H} = \sum_{i=1}^{N} \left( -\frac{\hbar^2}{2m} \frac{d^2}{dx_i^2} + v(x_i) \right) + \sum_{i_1, i_2} g \delta(x_{i_1} - x_{i_2}), \quad (1)
\]

where \( \hbar \) is the reduced Planck constant, \( m \) is the atomic mass, and the total particle number is \( N = N_\uparrow + N_\downarrow \), where \( N_\uparrow \) and \( N_\downarrow \) are the number of particles of the two components (hereafter referred to as spin-up and spin-down particles). The index \( i = 1, \ldots, N \) labels all particles (irrespectively of their spin state), while the indices \( i_1 \) and \( i_2 \) label, respectively, only spin-up and only spin-down particles. We focus on a balanced (unpolarized) mixture of the two components \( N_\uparrow = N_\downarrow = N/2 \). The external potential \( v(x) = V \sin^2(\pi x/d) \) represents the effect of an optical lattice with period \( d \) and intensity \( V \). The latter will be conveniently expressed in units of the recoil energy \( E_r = \hbar^2 \pi^2/(2md^2) \). We focus on a half-filled lattice, where the average density is \( n = N/(Ld) = 1/d \). The linear system size is \( Ld \), being \( L \) the number of wells of the OL. This is consistent with the use of periodic boundary conditions. The interaction strength is fixed by the one-dimensional coupling constant \( g = -2\hbar^2/(ma_{1D}) \), where \( a_{1D} \) is the one-dimensional scattering length. We consider the case of repulsive interactions \( g \geq 0 \). In the experiments preformed with atomic clouds confined in tight cigar-shaped waveguides, the coupling constant \( g \) can be related to the relevant experimental parameters [25], namely the three-dimensional s-wave scattering length and the radial harmonic confining frequency (assumed to be sufficiently strong to freeze the radial modes). Following the conventional formalism of homogeneous one-dimensional Fermi gases [26], we cast the interaction parameter in the adimensional form \( \gamma = 2/(n|a_{1D}|) \).

The ground state properties of the Hamiltonian (1) are determined via quantum Monte Carlo (QMC) simulations based on the DMC algorithm [28]. While in generic many-fermion systems the sign problem hinders exact QMC simulations, in one dimension this pathology can be circumvented without introducing...
any systematic approximation since the exact nodal structure is known [29-32]. In order to reduce the statistical fluctuations we employ Jastrow-Slater trial wave functions. The details of our implementation of the DMC algorithm have been reported in Refs. [21, 24].

In order to compute unbiased expectation values of operators that do not commute with the Hamiltonian, we employ forward walking technique [32].

In Fig. 4 we report the interaction energy per particle $E_{int}/N = (E - E_{\gamma=0})/N$, where $E$ is the total energy of an interacting and a noninteracting ($\gamma = 0$) gas in an OL, respectively. These data correspond to the particle number $N = 26$. In fact, by performing a finite-size scaling analysis using particle numbers in the range $18 \leq N \leq 54$, we verified that with $N = 26$ the relative error due to the finite system size is below the statistical errorbars in the weak interaction regime, and still below 0.3% in the regime of strong interactions $\gamma \approx 3$. In the weak OL limit $V \to 0$, the DMC results converge to the equation of state for a homogenous Fermi gas. This was determined in Ref. [24] by numerically solving the set of integral equations obtained by Yang [2] via the Bethe Ansatz technique. $E_{int}$ increases with the interaction strength $\gamma$, but it saturates in the Tongs-Girardeau limit $\gamma \to \infty$, where the energy of a fully polarized gas in an OL, respectively. These data correspond to the particle number $N = 26$. In fact, by performing a finite-size scaling analysis using particle numbers in the range $18 \leq N \leq 54$, we verified that with $N = 26$ the relative error due to the finite system size is below the statistical errorbars in the weak interaction regime, and still below 0.3% in the regime of strong interactions $\gamma \approx 3$. In the weak OL limit $V \to 0$, the DMC results converge to the equation of state for a homogenous Fermi gas. This was determined in Ref. 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gas is reached. While in a shallow OL ($V \approx E_r$) this saturation occurs only for strong interactions $\gamma \gg 1$, in a deep OL ($V \gg E_r$) it occurs already for intermediate interaction strengths $\gamma \approx 1$, meaning that correlation effects are enhanced in deep OLs compared to shallow OLs.

In the deep OL limit, one expects that higher bands become irrelevant if the interaction strength is not strong enough to promote interband transitions. By expanding the field operator in the basis of (maximally localized) Wannier functions, removing higher-band contributions, and neglecting also beyond-nearest-neighbor and interaction-induced processes (e.g., bond-charge interaction), the Hamiltonian \( H \) can be mapped to a discrete single-band lattice model, namely the (one-dimensional) Hubbard model \( \hat{H} = -t \sum_{\sigma} (\hat{c}^{\dagger}_{r,\sigma} \hat{c}_{r+1,\sigma} + h.c.) + U \sum_r \hat{n}_{r,\uparrow} \hat{n}_{r,\downarrow} \),

\begin{equation}
\hat{H} = -t \sum_{r,\sigma} (\hat{c}^{\dagger}_{r,\sigma} \hat{c}_{r+1,\sigma} + h.c.) + U \sum_r \hat{n}_{r,\uparrow} \hat{n}_{r,\downarrow},
\end{equation}

where \( \hat{c}^{\dagger}_{r,\sigma} (\hat{c}_{r,\sigma}) \) creates (destroys) a fermion of spin \( \sigma \in \uparrow, \downarrow \) at site \( r \) (with \( r = 1, \ldots, L \)), and \( \hat{n}_{r,\sigma} = \hat{c}^{\dagger}_{r,\sigma} \hat{c}_{r,\sigma} \) is the corresponding number operator. Consistently with the use of periodic boundary conditions, it is understood that \( \hat{c}^{\dagger}_{L+1,\sigma} = \hat{c}^{\dagger}_{1,\sigma} \) (\( \hat{c}_{L+1,\sigma} = \hat{c}_{1,\sigma} \)). The hopping energy \( t \) and the on-site interaction parameter \( U \) can be computed from Wannier functions integrals \( \sum_{\sigma} \) following the standard procedure \( 27 \).

The zero-temperature equation of state of the Hubbard model \( 22 \) was first determined by Lieb and Wu \( 3 \) using the Bethe Ansatz technique \( 35 \). The comparison displayed in Fig. \( 2 \) confirms that the continuous-space data do indeed converge to the Hubbard-model results if the OL is sufficiently deep. At the OL intensity \( V/E_r = 10 \), discrepancies are sizeable only for large values of the Hubbard interaction parameter \( U/t \gtrsim 10 \), which corresponds to the intermediate continuous-space interaction parameter \( \gamma \gtrsim 0.5 \).

Inducing strong-correlation effects in shallower lattices demands larger \( \gamma \) values. In this regime interband transitions become relevant; therefore, the mapping to a single-band model based on the standard Wannier function expansion \( 27 \) is invalid. As an attempt to take orbital excitations into account, we consider a parabolic approximation for the OL wells. The energy \( E_2 \) of two interacting opposite-spin fermions in the harmonic well can be exactly computed \( 38, 39 \). In the noninteracting case one has \( E_2 = \hbar \omega \), while in the Tonks-Girardeau \( (\gamma \to \infty) \) limit one has \( E_2 = 2 \hbar \omega \), as for two spin-aligned fermions. We henceforth define the one-site interaction parameter as the interaction energy \( U = E_2 - \hbar \omega \). In correspondence with the parabolic approximation for the interaction energy, we compute the hopping energy \( t \) using the well-known approximation - valid in the deep OL limit \( V/E_r \gg 1 \) - for (one fourth of) the bandwidth of the lowest band in the one-dimensional Mathieu equation, namely \( t = 4\pi^{-1/2} E_r (V/E_r)^{3/4} \exp \left(-2\sqrt{V/E_r}\right) \). This formula accounts to leading-order for the splitting of the harmonic-oscillator energy-levels due to tunneling \( 38, 39 \). The comparison of Fig. \( 2 \) shows that with this effective mapping procedure agreement between continuous-space and Hubbard-model data is obtained already at the moderate lattice depth \( V/E_r = 4 \), even when the continuous-space interaction parameter is as large as \( \gamma = 5 \) (where \( U/t \approx 11.45 \), according to this second mapping criterion).

Beyond the equation of state, we investigate how the antiferromagnetic correlations depend of the OL intensity and on the interactions strength. To quantify these correlations, we compute via DMC simulations the static magnetic structure factor of the continuous-space model: \( S_{\text{mag}}(k) = \langle \rho_{\text{mag}}(k) \rho_{\text{mag}}(-k) \rangle \), where \( \rho_{\text{mag}}(k) \) is the Fourier transform of the spin density operator. The results for \( S_{\text{mag}}(k) \) are shown in Fig. \( 3 \). In the upper panel, the different datasets correspond to different values of the interaction strength at the same OL intensity; in the lower panel, the OL intensity varies while the interaction strength is fixed. The peak of \( S_{\text{mag}}(k) \) at \( k = \pi/d \) signals antiferromagnetic correlations commensurate with the OL. One notices that such correlations can be amplified both by making the OL deeper and by increasing the interaction strength. In a deep OL of intensity \( V/E_r = 10 \), where the single-band description is applicable (see below), strong correlations emerge already at the moderate interaction strength \( \gamma \approx 0.2 \) (see Fig. \( 3 \)). However, even in OLs as shallow as \( V/E_r = 1 \), strong correlations form if the interaction parameter is close to the Tonks-Girardeau regime \( \gamma \gg 1 \), meaning that antiferromagnetism occurs also well beyond the tight-binding regime.

In order to analyze the convergence to the single-band limit, we make comparison with DMRG \( 40 \) results for the Hubbard model. Specifically, we compute via the DMRG method the spin-spin correlation function \( g(r_1, r_2) = \langle \hat{S}^z_r \hat{S}^z_{r_1} \rangle \), where the spin density operator is \( \hat{S}^z_r = \hat{n}_{r,\uparrow} - \hat{n}_{r,\downarrow} \). The Hubbard model results can be compared with the continuous-space magnetic structure factor using the following transformation (valid in the tight-binding limit) \( 38 \):

\begin{equation}
S_{\text{mag}}(k) = 1 + G^2(k) \left[ S_{\text{mag}}(k) - 1 \right],
\end{equation}

where \( S_{\text{mag}}(k) = 1/N \sum_{r_1, r_2} g(r_1, r_2) \exp[ik(r_1 - r_2)] \), and \( G(k) = \int |w_r(x)|^2 \exp(-ikx)dx \). It is worth stressing that in one dimension the spin density \( \langle \hat{S}^z_r \rangle \) is strictly zero as a consequence of the Mermin-Wagner theorem. The comparison between continuous-space and Hubbard model data is displayed in Fig. \( 4 \). We find that at the moderate OL depth \( V/E_r = 4.5 \) sizable discrepancies still persist, but in a deeper OL of intensity \( V/E_r = 10 \) precise matching is achieved. While the antiferromagnetic correlations are pronounced at large \( V/E_r \) and/or strong \( \gamma \), the Mermin-Wagner
finite arithmic divergence of the peak value order. This behavior implies, to leading order, a log-
technique [42, 43, 49], indicates quasi long-range spin corrections [50]. The value of the exponent
in the continuous-space notation), with logarithmic b

V/E f

fit

with system size, again well described by the logarithmic
short distance the spin-spin correlation functions of
long-range density-density correlations [49]. While at

density form a Wigner crystal characterized by quasi
continuous-space systems with (long-range) Coulomb repulsion (and no external potentials), which at low
density form a Wigner crystal characterized by quasi

limit. The system-size dependence of both the DMC
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