HST IMAGES REVEAL DRAMATIC CHANGES IN THE CORE OF IRC+10216

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ABSTRACT

IRC+10216 is the nearest carbon star with a very high mass-loss rate. The existence of a binary companion has been hinted at by indirect observational evidence, such as the bipolar morphology of its nebula and a spiral-like pattern in its circumstellar material; however, to date, no companion has been identified. We have examined archival Hubble Space Telescope images of IRC+10216, and find that the images taken in 2011 exhibit dramatic changes in its innermost region from those taken at earlier epochs. The scattered light is more spread out in 2011. After proper motion correction, the brightest peak in 2011 is close to, but not coincident with, the dominant peak in previous epochs. A fainter point-like object was revealed at ~0″.5 from this brightest peak. We suggest that these changes at the core of IRC+10216 are caused by dissipation of intervening circumstellar dust, as indicated by the brightening trend in the light curve extracted from the Catalina photometric survey. We tentatively identify the brightest peak in 2011 as the primary star of IRC+10216 and the fainter point-like source as a companion. The cause of non-detections of the companion candidate in earlier epochs is uncertain. These identifications need to be verified by monitoring the core of IRC+10216 at high resolution in the future.

Key words: circumstellar matter – stars: AGB and post-AGB – stars: individual (IRC+10216) – stars: late-type – stars: mass-loss – stars: winds, outflows

1. INTRODUCTION

Direct detection of binary companions of asymptotic giant branch (AGB) stars is challenging because the companions are overwhelmed by the luminous AGB stars. Moreover, companions of very dusty “infrared” AGB stars can be totally obscured by the circumstellar envelopes (CSEs). To date, binary companions have been detected for only a small number of AGB stars with low mass-loss rates through composite spectra or light curves (Jorissen et al. 2003). There are indications that a significant fraction of AGB stars have binary companions, such as elliptical CSEs (Mauron et al. 2013) as expected from binary AGB stars (Huggins et al. 2009), and spiral-shell patterns in the CSEs (Mauron & Huggins 2006; Mayer et al. 2011, 2014; Maercker et al. 2012; Kim et al. 2013; Cernicharo et al. 2013; Decin et al. 2015) modeled by AGB winds under the gravitational influence of binary companions (Theuns & Jorissen 1993; Soker 1994; Mastrodemos & Morris 1999; Kim & Taam 2012a, 2012b, 2012c; Kim et al. 2013).

Since its discovery (Becklin et al. 1969), IRC+10216, the nearest carbon star, has been the subject of multi-wavelength observations aimed at detecting new molecules or studying the circumstellar dust and chemistry (e.g., Guelin & Thaddeus 1977; Jura 1983; Glassgold 1996; Guelin 2011). The overall shape of its CSE, detected up to a radius of ~200″, is fairly spherical. Due to its proximity (distance ~130 pc; Groenewegen et al. 1998, 2012), the envelope has been resolved, exhibiting multiple shells with noticeably different geometric centers (Crabtree et al. 1987; Mauron & Huggins 1999, 2000; Fong et al. 2003; Decin et al. 2011). The dynamic timescale of shell separation, 200–800 years, is far longer than the pulsation period of IRC+10216 (~649 days; Le Bertre 1992) and significantly shorter than the periods between thermal pulses (of the order of 104 years; Bloecker 1995). The shell geometry must have been caused by other mechanisms, possibly binary interactions, as considered previously (e.g., Mauron & Huggins 1999; Decin et al. 2011). Indeed, the shell center offsets can be reproduced by a spiral-shell model of IRC+10216 with enhanced mass loss when a companion approaches periastron (e.g., Cernicharo et al. 2015). Recent Atacama Large Millimeter/submillimeter Array (ALMA) observations show that the morphology and kinematics of its inner parts are consistent with those expected from binary interactions (Decin et al. 2015).

High-resolution optical images of IRC+10216 taken in the years around 2000 revealed a bipolar nebula ~2″–3″ in size, with the northeast lobe significantly fainter than the southwest lobe (e.g., Haniff & Buscher 1998; Skinner et al. 1998; Mauron & Huggins 2000; Leão et al. 2006). This nebula was further modeled as an inclined bi-conical structure (Dyck et al. 1987; Haniff & Buscher 1998). Upon examination of all archival Hubble Space Telescope (HST) images of IRC+10216’s core, we find that in 2011 the aforementioned bipolar nebula had disappeared and a different set of peaks appeared. We identify the brightest peak and a fainter point-like source as the central star and its companion. In this Letter, we report our analysis of archival data of IRC+10216 and discuss its physical nature.

2. ARCHIVAL DATA

2.1. HST Images

Three images of IRC+10216 taken with the F606W filter were retrieved from the HST archives: Epoch 1 on 1998 March 30 with the WFC2 camera (PID 6856; PI: John Trauger) for a total integration time of t = 4000 s; Epoch 2 on 2001 January 7 with the WFC2 and t = 600 s (PID 8601; PI: Patrick Seitzer); and Epoch 3 on 2011 June 4 with the WFC3 camera and t = 5407 s (PID 12205; PI: Toshiya Ueta). The WFC2 and

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WFC3 fields are about 2.5′ × 2.5′ and 3′ × 3′, and the pixel sizes are 0.1″ and 0.04″, respectively. The three F606W images are displayed in Figure 1. It is immediately apparent that a very drastic change of structure has occurred between Epochs 2 and 3. A F814W image taken at Epoch 3 (Figure 1(d)) from the same program is saturated at the brightest peak but otherwise shows an almost identical morphology to the F606W image.

To compare these images, an astrometric analysis was performed using the CCMAP package of IRAF (Tody et al. 1993). Three well-detected stars located within 40″ from IRC+10216 and identified in the Sloan Digital Sky Survey (SDSS) catalog (Aihara et al. 2011) were chosen for reference. Only one of these stars is well detected in the shallow Epoch 2 image; therefore, we use two compact galaxies and one star to bootstrap the astrometry from the Epoch 1 and 3 images. The combined uncertainty in the astrometric alignment between each pair of images is ≤ 0.08.

2.2. Photometric Monitoring Data from the Catalina Sky Survey

We used the Catalina Sky Survey (Drake et al. 2014) data to investigate the variability of IRC+10216. The Catalina light
and ignoring the linear term amplitude of 1.8 ± 0.2 mag yr$^{-1}$, respectively, with the uncertainties of 0.2 mag yr$^{-1}$.

The average pulsation period of best curve of IRC+10216, shown in Figure 2.

Figure 2. Optical and infrared light curves of IRC+10216, and fitted sinusoids. (a) Top panel: Catalina light curve for nine years (2005–2013). One SDSS $i$-band data point (2005 December) is plotted as a cyan circle (see the text for details). The fit (solid curve) consists of a sinusoidal term with a period of 640 ± 0.9 days and an amplitude of 1.8 ± 0.02 mag (dotted-dashed), and a linear term for the brightening trend of −0.16 ± 0.004 mag yr$^{-1}$ (dashed). Vertical lines indicate the three relevant HST epochs in Figure 1. (b) Bottom left panel: $HJKL'M$ light curves of Le Bertre (1992) spanning four years (1985–1989). Arbitrary vertical shifts were used for comparison, as noted in the legend. The average pulsation period of best fits is 627 ± 26 days, which is slightly, but not significantly, smaller than the model ignoring the linear term (649 days; Le Bertre 1992). The $J$, $H$, and $K$ light curves show systematic brightening trends at rates of $−0.3$, $−0.2$, and $−0.1$ mag yr$^{-1}$, respectively, with the uncertainties of 0.2 mag yr$^{-1}$. (c) Bottom right panel: $HJKL'M$ light curves of Shenavrin et al. (2011) for nine years (2000–2008). The line colors and vertical shifts for bands are the same as in (b). The average amplitude ($1.7 ± 0.2$ mag) and period (633 ± 6 days) are similar to, albeit slightly smaller than, the ones in (b), but the brightening trends in the near-infrared are not found. The pulsation phase matches the Catalina ephemeris.

3. OPTICAL AND NEAR-INFRARED VARIATIONS

3.1. Morphological Variations in the Core

The HST F606W images from the three epochs are shown in Figure 1. The 1998 and 2001 images are qualitatively similar, exhibiting a bright southwest lobe and a fainter northeast lobe, previously identified as a bipolar structure (e.g., Mauron & Huggins 2000). The 2011 image shows a dramatically different morphology: the brightest peak appears to be at a different location, and the region within 1″ from the peak has a more complex structure, including a point-like source at 0.25 east of the peak. There is no trace of the previously identified bipolar structure. Using a 6″ radius source aperture and a background annulus of radii 14″–16″, we measured the HST F606W magnitudes of IRC+10216 to be 17.29, 17.15, and 16.93 mag in 1998, 2001, and 2011, respectively.

Images from the three epochs have been aligned astrometrically with respect to background stars and galaxies; however, to properly determine the differences among the images, the proper motion of IRC+10216 has to be considered. Using radio
continuum observations made with the Very Large Array (VLA), the proper motion of IRC+10216 has been reported to be $35 \pm 1$ mas yr$^{-1}$ in R.A. and $12 \pm 1$ mas yr$^{-1}$ in decl. (Menten et al. 2012). For this proper motion, the brightest peak in 2011 would correspond to positions close to, albeit not exactly on, the peaks in 1998 and 2001. The offsets (both about 0′18) are small, but are larger than the uncertainty in the differential astrometry (±0.08). The observed peak positions from 1998, 2001, and 2011 and the expected positions of the 2011 peak are plotted in Figure 3.

It is interesting to compare the 2011 peak position with the central position of IRC+10216 reported from high-resolution radio observations. The radio continuum emission of IRC+10216 originates from its radio photosphere (Menten et al. 2006), and its J2000 coordinates observed with the VLA in epoch 2006.16 are 09h47m57s.4255, Decin et al. 2015). Comparing these positions with the optical position of the 2011 peak (09h47m57s.4189, +13°16′43″859; see Figure 3), we see a 0′3 offset between the optical peak in 2011 and the VLA radio peak expected in 2011. Considering that the absolute astrometry of HST can reach 0′3–0′5 when a large number of reference stars are available (Koekemoer et al. 2006), the uncertainty of our absolute astrometry of IRC+10216 based on only three reference stars must be at least 0′5. Therefore, we suggest that the radio peak and the 2011 optical peak are coincident within the uncertainty of the HST astrometry.

3.2. Trends in the Light Curves of IRC+10216

The Catalina light curve of IRC+10216 in Figure 2(a) shows a gradual trend of brightening from 2005 to 2014, in addition to the stellar pulsation of the carbon star. To quantify the brightening, we fit the light curve in Figure 2(a) with a sinusoidal and a linear term:

$$m(t) = (A_0/2)\cos(2\pi(t - A_1)/A_2) + (A_3 + A_4t),$$

where $m(t)$ is the observed magnitude at epoch $t$ and $A_0$, $A_1$, $A_2$, $A_3$, and $A_4$ are constants. The best-fit model has $A_1 = 1.8 \pm 0.02$ mag for the amplitude of the sinusoidal variations, $A_2 = 640 \pm 0.9$ days for the period of the sinusoidal variations, and $A_4 = -0.16 \pm 0.004$ mag yr$^{-1}$ for the brightening rate in the Catalina photometry. The derived pulsation period of the carbon star is consistent with the range of values from the infrared bands (649 days, Le Bertre 1992; 630 days, Menten et al. 2012). No similar monitoring records at earlier epochs in the unfiltered or $i$ band are available to show the onset of the brightening trend. The $HST$ F606W measurements span 14 years but do not provide sufficient data points to constrain the parameters in Equation (1).

IRC+10216 was monitored in the infrared bands in 1985–1989 (Le Bertre 1992) and 2000–2008 (Shenavrin et al. 2011). These light curves are plotted in Figures 2(b) and (c) and fitted by Equation (1). We find that the $J$, $H$, and $K$ light curves in the 1980s show systematic brightening trends at rates of $A_4 = -0.3$, $-0.2$, and $-0.1$ mag yr$^{-1}$, but at longer wavelengths, the $L$ and $M$ light curves have nearly null brightening rates. Interestingly, the brightening trends in the near-infrared wavelengths vanished in the 2000s, when all observed infrared magnitudes stayed at a constant level with smaller amplitudes of sinusoidal variations (see Figure 2(c)). The $H$, $K$, $L$, and $M$ magnitudes seem to become fainter by 0.3–0.7 mag from epochs 1989.5 to 1999.9.

The light curves of IRC+10216 are clearly complex. The important features that need to be noted are that (1) there is the brightening trend of broad-band optical magnitudes in 2005–2014 and of near-infrared bands in 1985–1989; (2) the near-infrared $JHK$ magnitudes vary around nearly constant levels in 2000–2008, although the optical magnitudes show brightening in 2005–2008; and (3) the mid-infrared bands of $L$ and $M$ show periodic variations around constant levels during all monitoring campaigns, regardless of the brightening trends in other bands.

4. DISCUSSION

4.1. Nature of the Changes in the Core Region

The innermost region in the CSE of IRC+10216 in the 2011 $HST$ image exhibits a strikingly different appearance from those of 1998 and 2001 (Figure 1). This large modification of the optical core structure has occurred in $\lesssim$10 years; optical and infrared light curves have both shown brightening trends with a similar timescale. This timescale is much shorter than the stellar evolutionary timescale ($\sim$$10^3$–$10^4$ years) associated with thermal pulsing behavior and is unlikely due to variation in the stellar luminosity. Since the dust envelope of IRC+10216 is optically thick, the stellar radiation is absorbed and reemitted mostly in infrared wavelengths; its SED peak is near 10 $\mu$m or beyond in the mid-infrared (Ladjal et al. 2010). The lack of
long-term variations in the mid-infrared $L$ and $M$ bands indicates that the stellar luminosity is steady.

The timescale of $\lesssim 10$ years may be appropriate for photospheric or inner circumstellar changes. For example, a progressive rarefaction of the inner envelope allows penetration of stellar radiation, which can be detected in optical and near-infrared bands as a gradual brightening. The brightening trend could also be variable scattering due to changes in the geometry of the circumstellar material. Therefore, the brightening trends in the optical and near-infrared light curves likely imply the change of intervening circumstellar dusts in the $\lesssim 10$ year timescale, which may account for the recent changes of optical morphology in the core of IRC+10216.

4.2. Optical Identification of Central Star

The 2011 HST F606W image of IRC+10216’s core shows an elongated feature centered on the brightest peak. We established in Section 3.1 that this peak is coincident with the expected radio peak within the uncertainties of absolute astrometry; therefore, we identify this brightest peak as the central star of IRC+10216. The expected locations of this central star in 1998 and 2001 are both offset from the brightest peak by $\sim 0\prime.18$ (or 23 AU). This offset is larger than the uncertainty in the relative astrometry ($\lesssim 0\prime.08$). It is likely that the brightest peaks in 1998 and 2001 were both dust clumps in the vicinity of the central star.

A dust lane of IRC+10216, extending up to 4" in the east–west direction, is recently presented by Jeffers et al. (2014) in a linearly polarized optical intensity map. By placing the brightest peak in the HST’2011 image at their suggested location of the central star in the dust lane, we noticed that our elongated structure (P.A. $\sim -15^\circ$) is perpendicular to their dust lane (P.A. $\sim 80^\circ$) and connects the brightest parts in the two lobes of polarized optical light. The width of our elongated structure is $\sim 10$–20 AU (or 5–10 stellar radii with the photospheric radius of IRC+10216 $\sim 2$ AU suggested by Menten et al., 2012), which is similar to the radius of a dust formation zone or the inner radii used in the torus and ring models of Jeffers et al. (2014). The elongated structure could be caused by either an illumination or a confined outflow through the central hole of a dense equatorial torus.

The elongated structure consists of two compact blobs on two opposite sides of our suggested central star. Both blobs are $\sim 0\prime.2$ from the central star; assuming the same expansion speed as the outer envelope ($\sim 15 \text{ km s}^{-1}$; Loup et al. 1993), they could have been launched about eight years ago. We also note that the elongated structure extends up to $\sim 0\prime.4$, which corresponds to about 26 stellar radii. The elongated structure is slightly curved, thus shaped like a spiral or a point-symmetric structure.

4.3. Candidate for Companion of IRC+10216

One of the exciting findings from the 2011 HST image is the emergence of a point-like source at 05 east of the central star (indicated by an arrow in Figure 1(d)). This source is projected within the dust lane of Jeffers et al. (2014). The FWHM of its profile, 0.6″, is within the range of 0″09–0″17 measured for 10 unsaturated images of isolated stars. The color ($I_{F814W}/F606W$) of this point-like source is redder than those of diffuse clumps scattered over the inner 1″ region (denoted by 1–4 in Figure 1(c)) by a factor of 2.9 ± 1.1, suggesting a different nature from the ambient matter. We posit that this point-like source is a companion star of IRC+10216.

Its $r$ magnitude, determined by referring to SDSS stars, is $r = 21.1 \pm 0.2$ mag. Assuming that circumstellar extinction follows the interstellar extinction law and adopting the distance 130 pc, one can derive physical parameters of this companion candidate. For an $r^{-2}$ density distribution of the envelope, the column density to the source at radius $R$ is calculated to be $0.5\pi R_d/R$ of that to the center, where $R_d$ represents the radius of the dust formation zone (15 AU is adopted). For a visual optical depth $\tau_{0.55,\text{mag}} = 25$ derived from an SED fitting by Ladjal et al. (2010), the visual extinction toward the companion candidate would be $A_V \approx 8$ mag. With this extinction, its $r$ magnitude best matches a main-sequence star with a spectral type of M1 ($\sim 0.5 M_\odot$). It is, however, very tentative because of uncertainties in the circumstellar extinction coefficients, inhomogeneity of the AGB wind, and an unknown line-of-sight location of the point-like source.

If the point-like source seen in 2011 is indeed a binary companion, we can assess its possible orbit and proper motion. Assuming that the projected separation of 0″5 (i.e., 65 AU) corresponds to the orbital semimajor axis $a$, the orbital velocity of a binary component with respect to the other is $v = 6 \text{ km s}^{-1} (M/2.5 M_\odot)^{1/2}$, where $M$ is the total binary mass. This speed corresponds to 0″01 yr$^{-1}$ for an assumed total binary mass of $2.5 M_\odot$. Assuming a fitting accuracy of $\sim 0″02$, its movement may be observable a few years after 2011 in the best-case scenario. The formula for binary dynamics, $(p/\text{years})^2 = (a/\text{AU})^3 (M/M_\odot)^{-1}$, yields the orbital period $p$ of 330 years ($M/2.5 M_\odot)^{1/2}$. This period is consistent with the timescale of 200–800 years derived from the intervals between the outer shell patterns (Craib et al. 1987; Mauron & Huggins 1999, 2000), which, in binary-induced spiral-shell models, represents the orbital period (e.g., Kim & Taam 2012b).

The candidate for companion of IRC+10216 is proposed based on its point-like appearance, projection onto the dust lane, and color difference from nearby diffuse clumps. Its photometric and orbital period properties are not inconsistent with a low-mass main-sequence star. However, the absence of the companion candidate in the earlier epochs remains an open question since its orbital motion, if circular, is rather slow to explain its relative position. A speed $>30 \text{ km s}^{-1}$ is required to move 0″5 in 10 years. From single-imaging data, we cannot rule out the possibility that this point-like source could be a compact clump originated from small-scale fluctuations in the dust condensation radius of IRC+10216. Further observations and quantitative modeling are desired to better investigate this object in the current and near future, especially since the giant star in IRC+10216 will reach the next light minimum in 2016 March when the companion is less overwhelmed.

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