The diameter and evolutionary state of Procyon A

Multi-technique modeling using asteroseismic and interferometric constraints

P. Kervella¹, F. Thévenin², P. Morel², G. Berthomieu², P. Bordé³ and J. Provost²

1 European Southern Observatory, Casilla 19001, Vitacura, Santiago 19, Chile
2 Département Cassini, UMR CNRS 6529, Observatoire de la Côte d’Azur, BP 4229, 06304 Nice Cedex 4, France
3 LESIA, Observatoire de Paris-Meudon, 5, place Jules Janssen, F-92195 Meudon Cedex, France

Received ; Accepted

Abstract. We report the angular diameter measurement obtained with the VINCI/VLTI instrument on the nearby star Procyon A (α CMi A, F5 IV-V), at a relative precision of ±0.9%. We obtain a uniform disk angular diameter in the K band of θGD = 5.376 ± 0.047 mas and a limb darkened value of θLD = 5.448 ± 0.053 mas. Together with the HIPPARCOS parallax, this gives a linear diameter of 2.048 ± 0.025 D⊙. We use this result in combination with spectroscopic, photometric and asteroseismic constraints to model this star with the CESAM code. One set of modeling parameters that reproduces the observations within their error bars are an age of 2.314 Myr, an initial helium mass fraction Y1 = 0.301 and an initial mass ratio of heavy elements to hydrogen (Z⊙) = 0.0314. We also computed the adiabatic oscillation spectrum of our model of Procyon A, giving a mean large frequency separation of ∆ν0 ≈ 54.7 μHz. This value is in agreement with the seismic observations by Martin et al. [1999, 2001]. The interferometric diameter and the asteroseismic large frequency spacing together suggest a mass closer to 1.4 M⊙ than to 1.5 M⊙. We conclude that Procyon is currently ending its life on the main sequence, as its luminosity class indicates.

Key words. Stars: individual: α CMi, Stars: fundamental parameters, Stars: evolution, Techniques: interferometric

1. Introduction

Procyon A (α CMi, HD 61421, HIP 37279) is among the brightest stars in the sky and is easily visible to the naked eye. This made it an ideal target for a number of spectro-photometric calibration works. It is also a visual binary star (ADS 6251A) classified F5 IV-V, with a white dwarf (WD) companion orbiting the main component in 40 years. The influence of this massive companion on the apparent motion of Procyon was discovered by Bessel [1844]. Girard et al. [2000] have measured precisely the orbit of the pair, and obtained masses of 1.497 ± 0.037 M⊙ and 0.602 ± 0.015 M⊙, respectively for Procyon A and B. It has also been an asteroseismic target since a decade and Martin et al. [1999] and [2001] (hereafter M99 and M01) have measured a large frequency spacing of respectively 55 and 54 μHz. These asteroseismic observations provide strong constraints on stellar interior models, and on macroscopic stellar parameters like mass and radius. Comparing the direct interferometric measurements of this star to its model diameter is therefore important to cross-validate both approaches.

Matching Procyon’s position in the HR diagram has been recognized as of great difficulty by Guenther & Demarque [1993] (hereafter GD93). The reason is the poorly known mass and metallicity Z of the star. We note that among the computed models only their model b is close to the mean large frequency spacing measured by M99-M01. Because numerous new studies and observational constraints (like the direct diameter) exist today, we re-examine the status of Procyon A in helium content Y and in age. We first review and present the adopted fundamental parameters of Procyon A in Sect. 2 part of them being used for the limb-darkening in Sect. 3 where we present our new interferometric observations and the associated data processing. In Sect. 4 we detail several models of Procyon computed with the CESAM code (Morel [1997]). These models are constrained using the spectroscopic effective temperature and the linear diameter value that we derived from the VINCI/VLTI observations. Thus, we avoid to fit the luminosity for which bolometric corrections are rather uncertain for stellar luminosity class IV-V. A study of the asteroseismic frequencies is then presented in Sect. 5 before the conclusion.
2. Fundamental parameters

Physical properties of Procyon A are well known thanks to the carefully measured orbit of the system and accurate 
HIPPARCOS parallax. Its mass has been recently measured 
by Girard et al. (2000) at \( M_A = 1.497 \pm 0.037 M_\odot \), adopting 
a parallax of 283.2 \pm 1.5 mas. This mass is computed 
using a large number of observations of inhomogeneous 
quality. The same authors when they limit their observation 
sample to the excellent images obtained with the 
WFPC2 found \( M_A = 1.465 \pm 0.041 M_\odot \). If we replace 
the parallax used by Girard et al. (2000) by the parallax from 
the HIPPARCOS catalogue, 285.93 \pm 0.88 mas, the sum of 
the masses of the binary decreases by 2.9 \%. Consequently, 
masses become respectively \( M_A = 1.42 \pm 0.04 M_\odot \) for 
Procyon A and \( M_B = 0.575 \pm 0.017 M_\odot \) for the white 
dwarf Procyon B. Allende Prieto et al. (2002) hereafter 
AP02, using the HIPPARCOS parallax of Procyon A, have 
estimated its mass to be 1.42\pm0.06\,M_{\odot}. From these works, 
we can conclude that the mass of Procyon is probably 
between 1.4 and 1.5\,M_{\odot}. Consequently, we investigate 
these two values in our study. The other adopted stellar param-
eters are summarized in Table 1.

The bolometric luminosity is given by Steffen (1985)
\[ \log (L_b / L_\odot) = 0.85 \pm 0.06. \]
The photometric properties of Procyon have been carefully studied by AP02, who 
derived an effective temperature \( T_{\text{eff}} = 6512 \, K \). This value is 
based on the angular diameter measured in the visible by 
the MarkIII optical interferometer (Mozurkewich et al. 
1991). Detailed 3D stellar atmosphere simulations have 
led AP02 to correct this value of \( T_{\text{eff}} \) by 80 K. Following 
their results, we adopt \( T_{\text{eff}} = 6530 \pm 50 \, K \) in our com-
putations with a surface gravity of \( \log g = 3.96 \pm 0.02. \) 
It is interesting to notice that over the last twenty years, 
the effective temperature estimates for Procyon are spread 
between 6545 and 6811 K (Cayrel et al. 1997). 
Adopting a color index of \( B - V = 0.421 \) and using the calibration by 
Alonso (1990) we find \( T_{\text{eff}} = 6551 \, K \) which is very close to 
the value proposed by AP02.

The iron abundance at surface of [Fe/H] = \(-0.05\pm0.03\) 
dex with respect to the solar one is also taken from AP02. 
Therefore we adopt for the modeling of the internal structure 
of Procyon A a chemical mixture which is calculated from 
the iron abundance using the approximation:

\[
\text{[Fe/H]}_s \equiv \log \left( \frac{Z_{\text{Fe}}}{Z} \right)_s + \log \left( \frac{Z}{X} \right)_s - \log \left( \frac{Z_{\text{Fe}}}{X} \right)_\odot \tag{1}
\]

\[
\simeq \log \left( \frac{Z}{X} \right)_s - \log \left( \frac{Z}{X} \right)_\odot
\]

where \( \left( \frac{Z_{\text{Fe}}}{Z} \right)_s \) is the iron mass fraction within Z. We use 
the solar mixture of Grevesse & Noels (1993) then \( \left( \frac{Z}{X} \right)_\odot = 0.0245 \). 
As we shall see in Sect. 3, we adopt at the age of 
Procyon A a surface abundance \( \left( \frac{Z}{X} \right)_s = 0.0217 \pm 0.0017 \). 
Note that this adopted error bar of AP02 is very small. 
Because the white dwarf (WD) Procyon B has already 
experienced the AGB phase, it is not excluded that material 
coming from this star can have contaminated the narrow 
surface convection layer of Procyon A. However, there is no 
evidence of such an effect in the published table of abund-
ces of Procyon A (Steffen 1985). Enriched elements like 
“s” or other elements coming from Procyon B during its 
post-AGB phase do not seem to have migrated as gas or 
dust in the atmosphere of Procyon A. The mass transfer 
during this short period of the WD progenitor life could 
not exceed the accretion of the wind passing close to the 
star. Therefore, except if the two stars have filled in their 
Roche lobe, the mass of Procyon A has only been changed 
by a negligible amount.

The projected rotational velocity of Procyon A has 
been determined by many authors who derived values 
smaller than \( 5.0 \, \text{km}\,\text{s}^{-1} \). Considering in particular 
the value proposed by AP02, \( v \sin i = 3.2 \pm 0.5 \, \text{km}\,\text{s}^{-1} \), we 
conclude that Procyon A is a slow rotator and we 
neglect the rotational velocity in our modeling of its internal 
structure. The semi-major axis of the Procyon orbit being 
\( \alpha \simeq 4''5 \) (Girard et al. 2000), the distance between the 
two companions amounts to \( \simeq 16 \, \text{au} \). The tidal interaction 
between the two components is therefore negligible.

3. Interferometric angular diameter

3.1. Instrumental setup and data processing

The European Southern Observatory’s Very Large 
Telescope Interferometer (Glindeemann et al. 2000) 
is operated on top of the Cerro Paranal, in Northern Chile, 
since March 2001. The observations reported here were 
done with two test siderostats (0.35 m aperture) and the 
VINCI beam combiner (Kervella et al. 2000, 2003a). 
The Procyon visibility measurements were all obtained on 
the B3-D3 baseline, 24 m in ground length. We used a regular 
K band filter (\( \lambda = 2.0 - 2.4 \, \mu m \)) for these observations. 
The effective wavelength changes slightly depending on 
the spectral type of the observed target. For Procyon, we 
estimate \( \lambda_{\text{eff}} = 2.178 \pm 0.003 \, \mu m \). This error bar adds 
quadratically a relative systematic uncertainty of 0.15\% 
on the final limb darkened angular diameter error bar.

The raw data processing has been achieved using an 
improved version of the standard data reduction pipeline 
(Kervella et al. 2003b), whose general principle is similar 
to the original FLUOR algorithm (Coudé du Foresto et al. 
1997). The two calibrated output interferograms are 
subtracted to remove residual photometric fluctuations. 
Instead of the classical Fourier analysis, we implemented 
a wavelet based time-frequency analysis (Ségransan et al. 
1999). The output of the processing pipeline is a single 
value of the squared coherence factor \( \mu^2 \) for each series of 
500 interferograms and the bootstrapped error bar.

3.2. Visibility calibration

Sirius was chosen as the calibrator for Procyon obser-
vations as its limb darkened (LD) angular diameter has 
been measured recently with high precision by Kervella 
et al. (2003b), using VINCI and the 66m E0-G1 base-
Table 1. Relevant parameters of α CMi (Procyon) and its calibrator α CMa (Sirius). For Sirius, see also Kervella et al. (2003d).

|         | α CMi       | α CMa       |
|---------|-------------|-------------|
| m_V     | 0.34        | -1.47       |
| m_K     | -0.65       | -1.31       |
| Spectral Type | F5IV-V      | A1Vm       |
| M(M☉)   | 1.42 ± 0.04 | 2.12 ± 0.02 |
| T_eff(K)| 6530 ± 50   | 9900 ± 200  |
| log g   | 3.96        | 4.30        |
| [Fe/H]  | -0.05±0.03  | 0.50±0.20   |
| v sin i | 3.16±0.50   | 16.0±1.0    |
| π^a (mas)| 285.93 ± 0.88 | 379.22 ± 1.58 |
| θ_{UD} (mas) | 5.376 ± 0.047 | 5.94 ± 0.02 \(^b\) |
| θ_{LD} (mas) | 5.448 ± 0.053 | 6.04 ± 0.02 \(^b\) |

\(^a\) Parallax values from HIPPARCOS (Perryman et al. 1997).
\(^b\) Sirius θ_{UD} and θ_{LD} are taken from Kervella et al. (2003d).

line. Table 1 gives the relevant parameters of Procyon and Sirius. The interferometric efficiency (IE) of the instrument was estimated from the ρ^2 values obtained on Sirius a short time before and after Procyon (typically 15 minutes). Thanks to the very high SNR of the Sirius interferograms, we were able to obtain very small relative errors on the IE values, typically 0.2% for a series of 500 interferograms. The error bar on the angular diameter of Sirius is treated as a systematic error in the calibration process, and is therefore not averaged for multiple observations. For the visibility fit, we took into account simultaneously the limb darkening and the bandwidth smearing, as described in Kervella et al. (2003b, 2003d).

Our February 2003 observation campaign of Procyon resulted in a total of 53 calibrated V^2 measurements, from 23 256 processed interferograms. In the calibration process, we separated clearly the statistical and systematic error contributions. This is essential to avoid underestimating the final error bars of the fit. The latter corresponds to the uncertainty on the angular diameter of Sirius. In spite of the fact that the angular size of Sirius is larger than Procyon’s, the small error on its measured value (±0.3%) translates into a small calibration error. As a consequence, the global error on the angular size of Procyon is dominated by the statistical contribution by a factor of two compared to the systematic calibration error: the uncertainty on the angular size of Sirius is not limiting our final precision. The visibility points as a function of the projected baseline are presented on Fig. 1 and their residual scattering around the best fit model on Fig. 2.

3.3. Limb darkened angular diameter and linear size

The intensity profile I(μ)/I(1) that we chose for Procyon was computed by Claret et al. (2000), based on the

\[ \frac{I(x)}{I(1)} = 1 - \sum_{k=1}^{4} a_k (1 - x^k) \]  

which coefficients are a_1 = 0.5089, a_2 = 0.0822, a_3 = -0.3978, and a_4 = 0.1967. The parameter x = cos(θ) is the cosine of the azimuth of a surface element of the star. The resulting profile is shown on Fig. 3. AP02 have identified an influence of the convective granulation of Procyon on its limb darkening at visible wavelengths. Though, the amplitude of this effect is already very small at λ = 1.0 µm, and we neglect it for our observations at λ = 2.2 µm.

We obtain directly the LD angular diameter of Procyon from a classical χ^2 minimization, θ_{LD} = 5.448 ± 0.053 mas. Using a simple uniform disk model, we find a value of θ_{UD} = 5.376 ± 0.047 mas.

We can compare the LD value with the previously published interferometric measurements, listed in Table 2. Our value of θ_{LD} is compatible with previous measurements, that were all obtained using visible wavelength observations. It is also in excellent agreement with the average
of all published values. From the VINCI/VLTI value of θ_{LD} and the HIPPARCOS parallax, we deduce a linear diameter of 2.048 ± 0.025 D⊙. Computing the average of all published θ_{LD} values, we obtain a linear radius of R_{⋆} = 2.047 ± 0.020 R⊙ for Procyon A, statistically identical to our result.

### 4. Modeling

The modeling of Procyon A is based on the matching of spectroscopic effective temperature, surface metallicity and interferometric radius by a “satisfactory” evolutionary model for a given mass. The constraint between R⋆ and T_{eff} within their error bars is illustrated in Fig. 4 by the hatched parallelogram, while the constraint between the T_{eff} and L⋆ is illustrated by the dashed rectangle. That emphasizes the advantage of the use of the measured R⋆ instead of the L⋆ which depends of photometric calibrations and bolometric corrections.

Table 3 shows the characteristic and initial modeling parameters of the star. The ordinary assumptions of stellar modeling are made, i.e. spherical symmetry, no rotation, no magnetic field and no mass loss. The relevant nuclear reaction rates are taken from the NACRE compilation by Angulo et al. (1999). The equation of state adopted is taken from Eggleton et al. (1978), and the opacities are from the OPAL database (Iglesias & Rogers 1996), using the Grevesse & Noels (1993) mixture. The microscopic diffusion is described using the formalism of Burgers (1969) with the resistance coefficients of Paquette et al. (1990).

We take into account the radiative diffusivity as recommended by Morel & Thévenin (2002), that limits the efficiency of the microscopic diffusion in outer-layers of stars with intermediate masses. The atmosphere is restored using Hopf’s law (Mihalas 1978). The definition of the radius of the star is the bolometric radius, where T(R⋆) = T_{eff}. In the convection zones the temperature gradient is computed according to the MLT_{CM} convection theory with a mixing length parameter of Λ = 1 (Canuto & Mazzitelli 1991, 1992).

Following the discussion of Sect. 2 we investigate the sensitivity of our models to a variation of the mass of Procyon in the range 1.4 to 1.5 M⊙. We obtained a “satisfactory” model a (Fig. 4), detailed in Table 3 which reaches the hatched parallelogram corresponding to an age of 2.314 Myr. Its evolutionary state corresponds closely to the end of the main sequence, when the convective core has disappeared owing to the exhaustion of the hydrogen at center. We present on Fig. 5 the variation for model a of X_{s}, Y_{s} and of [Fe/H]_{s} as a function of the age. From the ZAMS until about 1 Gyr, due to the gravitational settling, the surface abundances of helium and heavy elements (resp. hydrogen) decrease (resp. increases). After 1 Gyr, the density in the envelope decreases due to the enlargement of radius. This leads to an increase of the radiative mixing, with the consequence of a dredge-up which increases (resp. decreases) Y_{s} and [Fe/H]_{s} (resp. X_{s}) at the surface.

We also computed a “satisfactory” model b without microscopic diffusion in order to estimate its effect on the age of the star and we found that suppressing the diffusion increases the age by ~ 400 Myr. The evolutionary state of this model is the beginning of H burning in shell.

Following the prescriptions of Schaller et al. (1992) we have computed additional models that include overshooting of the convective core of radius R_{co}, over the distance O_{s} = A \min(H_{p}, R_{co}). It has been impossible to find a “satisfactory” model reaching the narrow hatched area of the HR diagram with parameters A > 0.03.

Finally, we computed a model c with Girard et al.’s (2000) mass : 1.50 M⊙ and find an age of 1.3 Gyr without overshoot. We also computed a model with an overshoot of 0.15 in agreement with Ribas, Jordi & Carne (2001) for stars of mass lower than 1.5 M⊙ resulting in a model which increases the age of the star to 1.55 Gyr.

The parameters for all models are listed in Table 3. Fig. 4 shows the corresponding evolutionary tracks in the HR diagram. We note that our models fit the value of log g = 3.96 proposed by AP02.

### 5. Asteroseismology

Asteroseismic observations of Procyon A have been reported by M99 and M01 and indicate an excess of power around 1 millihertz with a probable frequency spacing of 54 to 55 μHz.
The set of frequencies of model order 450 to 1 600 range, i.e. large radial orders, \( \Delta \nu \) mean value radius.

Domain for effective temperature and the interferometric temperature, while the hatched area delimits the uncertainty domain for luminosity and effective temperature and the interferometric radius.

Fig. 4. Evolutionary tracks in the HR diagram of models of Procyon A (see Table 3). Dashed rectangle delimits the uncertainty domain for luminosity and effective temperature, while the hatched area delimits the uncertainty domain for effective temperature and the interferometric radius.

Fig. 5. Changes with respect to time for \( X_s \), \( Y_s \) and \([Fe/H]_s\) for model a of Procyon.

The narrow convective envelope of the star we consider may stochastically excite solar-like oscillations. As in the Sun, these oscillations will have quasi equidistant frequencies at given degree.

For all the models, we have computed a set of adiabatic frequencies of the normal modes for degrees \( \ell = 0, 1, 2, 3 \). The set of frequencies of model a are given in the frequency range 450 to 1 600 \( \mu Hz \) in Table 3.

The large frequency spacing is defined as the difference between frequencies of modes with consecutive radial order \( n \): \( \Delta \nu_l(n) \equiv \nu_{n,\ell} - \nu_{n-1,\ell} \). In the high frequency range, i.e. large radial orders, \( \Delta \nu_l \) is almost constant with a mean value \( \Delta \nu_0 \), strongly related to the mean density of the star. The computed frequencies are fitted to the following asymptotic relation (Berthomieu et al. 1993):

\[
\nu_{n,\ell} = \nu_0 + \Delta \nu_0 (n - n_0) + a_\ell (n - n_0)^2. \]

With \( n_0 = 21, \ell = 0 \) and radial order \( n \) between 17 and 26 we obtain an estimate of the mean large difference \( \Delta \nu_0 \).

Table 3. Procyon A models (without overshoot, see text) lying within uncertainty box in the HR diagram. The subscripts “i” and “s” respectively refer to initial values and surface quantity at present day. “c” refers to the central value.

| Model | a | b | c |
|-------|---|---|---|
| \( m/M_\odot \) | 1.42 | 1.42 | 1.50 |
| \( Y_i \) | 0.3012 | 0.2580 | 0.345 |
| \( Y_0 \) | 0.2209 | 0.2580 | 0.202 |
| \( (\frac{V}{T_i})_s \) | 0.03140 | 0.0218 | 0.0450 |
| \( (\frac{V}{T_0})_s \) | 0.02157 | 0.0218 | 0.0220 |
| diffusion | yes | no | yes |
| age (Myr) | 2.314 | 2.710 | 1.300 |
| \( T_{\text{eff}} \) (K) | 6524 | 6547 | 6553 |
| log \( g \) | 3.960 | 3.967 | 3.994 |
| \( [Fe/H]_i \) | +0.107 | -0.051 | +0.264 |
| \( [Fe/H]_s \) | -0.055 | -0.051 | -0.043 |
| log(\( L/L_\odot \)) | 0.8409 | 0.8405 | 0.8390 |
| \( R/R_\odot \) | 1.300 | 0.107 | +0.264 |
| \( \Delta \nu_0 (\mu Hz) \) | 54.7 | 55.4 | 56.4 |

The mean large spacing of model a is 54.7 \( \mu Hz \) which fits well observational predictions of M99 & M01 (\( \Delta \nu_0 \sim 54 - 55 \mu Hz \)). For the model c, the mean large spacing is found to be 56.4 \( \mu Hz \), significantly different from the observed value. A mass as large as 1.5 \( M_\odot \) for Procyon A seems therefore improbable. Adopting the VINCI/VLTI radius, and looking at the scaling of Kjeldsen & Bedding (2003),

\[
\Delta \nu_{\text{osc}} \sim 134.9 \sqrt{\frac{m/M_\odot}{(R/R_\odot)^3}} (\mu Hz),
\]

the most straightforward way to decrease \( \Delta \nu_0 \) is to decrease the mass of Procyon A.

Because the luminosity \( L \) is in relation with \( T_{\text{eff}} \) and with the radius \( R \) through the formula \( L/L_\odot = (R/R_\odot)^2(T_{\text{eff}}/T_{\text{eff}})^4 \) we find from Eq. 3 the expression of the mean large spacing as a function of \( L \) and \( T_{\text{eff}} \),

\[
\Delta \nu_{\text{osc}} \sim 134.9 \sqrt{\frac{m/M_\odot}{(L/L_\odot)^3(T_{\text{eff}}/T_{\text{eff}})^4}} (\mu Hz). \]

From this formula it is easy, if needed, to estimate the derivatives of the mean large spacing with respect to the mass, the effective temperature and the luminosity. To estimate the dependance between \( \Delta \nu_{\text{osc}} \) and \( Z_\odot \), we have computed a modified model a with \( \Delta Z = 0.001 \) and found a variation in the parallelogram error box of: \( \Delta T_{\text{eff}} = 22 K, \Delta L/L = 0.008, \) and \( \Delta \alpha_{\text{age}} = 33 \) Myr. There were no significant change with \( Z_\odot \) of \( \Delta \nu_{\text{osc}} \) as can be expected from Eq. 3.

6. Discussion & Conclusions

We have reported in this paper our modelisation of Procyon A based on observations by long-baseline interferometry and asteroseismology. The use of the measured
stellar diameter allows to reduce significantly the error bar on the luminosity of the star, as it does not require any bolometric corrections. This advantage has been demonstrated in this study and in the modeling of Sirius by Kervella et al. (2003).

Our model shows that using the given set of parameters and given physics, Procyon A is currently finishing to burn its central hydrogen, and is at the phase where the convective core is disappearing. The error on the measured radius gives a narrow uncertainty of 10 Myr on the deduced age. Provencal et al. (2002) have discussed the cooling time of the WD Procyon B and found that the progenitor ended its lifetime 1.7 ± 0.1 Gyr ago. We derive an age of 2.314 Myr for Procyon A. Subtracting the cooling age of the WD companion to our determination of the age of Procyon A leads to a lifetime of 614 Myr for the progenitor of Procyon B. This indicates that the mass of the WD (see also Jeffries 1997). This estimate of 0.57 $M_\odot$ agrees very well with the mass of Procyon B that we deduced from Girard et al. (2001). We note that the age obtained with model $c$ is shorter than the cooling age of Procyon B. This argument suggests a mass lower than 1.5 $M_\odot$ for Procyon A and strengthens the asteroseismology results.

Further progress on the modeling of Procyon will be possible when the accuracy on the flux of the star is improved to less than 1%. Waiting for such accuracy, the uniqueness of the solutions resulting from computed models fitting a narrow box in the HR diagram will come from future detailed asteroseismic studies. For example, other linear combination of frequencies like the small frequency spacings (see e.g. Gough 1991) will constrain the age and the mass of the star. We therefore encourage observers to progress on the asteroseismology of this star.

Large uncertainties also come from the adopted chemical abundance mixture $Z_\odot$ which is still rather uncertain. Only a few chemical element abundances are measured today and most of them with a low accuracy. This uncertainty on $Z_\odot$ is the largest source of error on the estimated initial helium content $Y_i$ and on the age of Procyon. Thus, we recommend that surface abundances should be derived from 3D atmosphere studies, in particular for oxygen and other important donors of electrons.

Acknowledgements. The VLTI is operated by the European Southern Observatory at Cerro Paranal, Chile. The VINCI public commissioning data reported in this paper has been retrieved from the ESO/ST-ECF Archive. The VINCI pipeline includes the wavelets processing technique, developed by D. Ségransan (Obs. de Genève). No VLTI observation would have been possible without the efforts of the ESO VLTI team, for which we are grateful. This research has made use of the SIMBAD database at CDS, Strasbourg (France).

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Table 4. Asteroseismic frequencies (µHz) of Procyon A for model $a$. The “*” correspond to “g” modes which are presents in the frequency domain due to the evolutionary state of the star.

| $\ell=0$ | $\ell=1$ | $\ell=2$ | $\ell=3$ | $n$ |
|---------|---------|---------|---------|----|
| 485.35  | 513.36  | 532.31  | 551.46  | 8  |
| 537.93  | 563.43  | 585.90  | 605.61  | 9  |
| 591.09  | 615.78  | 638.60  | 657.72  | 10 |
| 643.23  | 666.79  | 689.80  | 708.98  | 11 |
| 694.17  | 717.44  | 741.00  | 760.87  | 12 |
| 745.32  | 769.15  | 793.75  | 814.22  | 13 |
| 798.11  | 822.24  | 847.62  | 868.87  | 14 |
| -       | -       | 896.38  | -       | *  |
| 852.14  | 876.76  | 904.48  | 923.72  | 15 |
| 906.98  | 931.22  | 957.38  | 978.12  | 16 |
| 961.42  | 985.26  | 1011.19 | 1031.92 | 17 |
| -       | -       | 1075.73 | -       | *  |
| 1015.27 | 1038.85 | 1065.03 | 1086.34 | 18 |
| 1069.16 | 1092.64 | 1119.00 | 1140.31 | 19 |
| 1123.16 | 1146.76 | 1173.64 | 1195.27 | 20 |
| 1177.86 | 1201.57 | 1228.79 | 1250.68 | 21 |
| 1233.05 | 1256.82 | 1284.20 | 1306.18 | 22 |
| 1288.46 | 1312.90 | 1339.65 | 1361.81 | 23 |
| 1343.91 | 1367.55 | 1395.21 | 1417.43 | 24 |
| 1399.42 | 1422.95 | 1450.65 | 1472.94 | 25 |
| 1454.83 | 1478.24 | 1506.00 | 1528.39 | 26 |
| 1510.12 | 1533.53 | 1561.40 | 1583.89 | 27 |
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