Lattice distortions, incommensurability, and stripes in the electronic model for high-T$_c$ cuprates

Takashi Yanagisawa, Mitake Miyazaki, Shigeru Koikegami, Soh Koike and Kunihiko Yamaji
Nanoelectronics Research Institute, National Institute of Advanced Industrial Science and Technology (AIST),
Central 2, 1-1-1 Umezono, Tsukuba, Ibaraki 305-8568, Japan
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Striped superconductivity (SC) with lattice distortions is investigated based on the three-band Hubbard model for high-T$_c$ cuprates. A stable inhomogeneous striped state is determined in the low-temperature tetragonal (LTT) phase with lattice distortions using a quantum variational Monte Carlo method. The ground state has vertical or horizontal hole-rich arrays coexisting with incommensurate magnetism and SC induced by several percents of lattice deformations. The SC order parameter oscillates according to the inhomogeneity in the antiferromagnetic background with its maximums in the hole-rich regions, and the SC condensation energy is reduced as the doping rate decreases.

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Over the last decade the oxide high-T$_c$ superconductors have been investigated intensively. The mechanism of superconductivity (SC) has been extensively studied using various two-dimensional (2D) models of electronic interactions. The 2D three-band Hubbard model is the simple and most fundamental model among such models. The 2D one-band Hubbard model is regarded as the simplified model of the three-band model. Studies of these models over the last decade indicated that the $d$-wave SC is induced from the electronic repulsive interaction; significantly it has been shown that the SC condensation energy and the magnitude of order parameter are in reasonable agreement with the experimental results in the optimally doped case.

The SC condensation energy obtained by the variational Monte Carlo method (VMC) is estimated as $E_{\text{cond}} \approx \frac{0.00117t}{t_{pd}} = 0.59\text{meV per site}$ in the optimally doped case for the single-band Hubbard model in the bulk limit. We must note that $E_{\text{cond}}$ is given as $0.17 \sim 0.26\text{meV}$ by specific heat data and $0.26\text{meV}$ by the critical magnetic field value $H_c^2/\pi$. The agreement of the VMC value with the experimental estimation is quite significant and supports the calculations. The VMC method can be regarded as an approximation to Quantum Monte Carlo calculations. The SC order parameter $\Delta_n$ determined from a minimum of the energy is of the order of $0.01\sim0.015\text{eV} \sim 20\text{meV}$ at hole density $\delta$ $\sim 0.2$.

The interplay between magnetism and superconductivity is suggested in the underdoped region. The reduction of T$_c$ in this region remains unresolved and may be related to magnetism. An existence of incommensurate correlations with modulation vectors given by $Q_s = (\pi, \pi, 0)$ and $Q_c = (\pi, \pi, 0)$ or $Q_s = (\pi, 2\pi, 0)$ has been suggested for the hole-doping rate $\delta$. The linear doping dependence of incommensurability in the underdoped region supports a striped structure and suggests a relationship between magnetism and SC. A relationship between the SDW, CDW orders and a crystal structure is also suggested in intensive studies by the neutron-scattering measurements. In particular in the low-temperature tetragonal (LTT) and low-temperature less-orthorhombic (LTLO) phases, the CDW order is stabilized, while no well defined incommensurate CDW peaks were observed for the orthorhombic systems. In the elastic and inelastic neutron scattering experiments with La$_{2-x}$Sr$_x$CuO$_4$, the incommensurate magnetic scattering spots around $(\pi, \pi)$ have been observed in the SC phase in the range of

![FIG. 1: SC condensation energy per site as a function of $1/N_a$ in $t$ units where $t = t_{pd}/3$ and $N_a$ is the number of atoms. Squares are for $\delta = 0.2$, $t_{pp}/t_{pd} = 0.4$ and $U/d/t_{pd} = 8$ for the three-band model on square lattices. Circles are at $\delta = 1/8$ coexisting with stripes for $t_{pp}/t_{pd} = 0.4$ and $U/d/t_{pd} = 8$ on rectangular lattices $32 \times 8$, $24 \times 6$, $16 \times 8$ and $16 \times 4$. Triangles are for the single-band Hubbard model; $\delta = 0.86$ and $t' = -0.2$ and $U = 8$ for solid symbols and $\delta = 0.84$ and $t' = -0.15$ for open symbols (energy unit is $t$). The diamond shows the value indicated from experiments.](image-url)
In this paper an incommensurate striped SC under lattice distortions is investigated based on the three-band model. In the description of SC in the underdoped region it is of prime importance to investigate the effect of inhomogeneity. In this paper we take into account all of the inhomogeneity, lattice instability and anisotropic pairing to clarify the ground state in the underdoped region of high-Tc cuprates. The Hamiltonian is given by 

$$H = H_{pd}^0 + V + H_{el}$$

where

$$H_{pd}^0 = \epsilon_d \sum_{i\sigma} d_{i\sigma}^\dagger d_{i\sigma} + \epsilon_p \sum_{i\sigma} [p_{i+\hat{x}/2,\sigma}^\dagger p_{i+\hat{x}/2,\sigma} + p_{i-\hat{y}/2,\sigma}^\dagger p_{i-\hat{y}/2,\sigma} - p_{i-\hat{x}/2,\sigma}^\dagger p_{i+\hat{y}/2,\sigma} + p_{i+\hat{x}/2,\sigma}^\dagger p_{i-\hat{y}/2,\sigma} + h.c.] + t_{pd} \sum_{i\sigma} [d_{i\sigma}^\dagger (p_{i+\hat{x}/2,\sigma} + p_{i-\hat{y}/2,\sigma}) - p_{i-\hat{x}/2,\sigma}^\dagger p_{i+\hat{y}/2,\sigma} + h.c.] + t_{pp} \sum_{i\sigma} [(1 + \mu_i) p_{i+\hat{x}/2,\sigma}^\dagger p_{i+\hat{y}/2,\sigma} - p_{i+\hat{y}/2,\sigma}^\dagger p_{i+\hat{x}/2,\sigma} - p_{i-\hat{x}/2,\sigma}^\dagger p_{i+\hat{y}/2,\sigma} + h.c.] + t_{pd} \sum_{i\sigma} [u_{i\hat{x}} d_{i\sigma}^\dagger p_{i+\hat{x}/2,\sigma} - u_{i\hat{y}} d_{i\sigma}^\dagger p_{i+\hat{y}/2,\sigma} + h.c.]$$

$$V = U_d \sum_i d_{i\uparrow}^\dagger d_{i\uparrow} d_{i\downarrow}^\dagger d_{i\downarrow},$$

where $H_{el}$ denotes the lattice elastic energy given by 

$$H_{el} = (K_{pd}/2) \sum_i (u_{i\hat{x}}^2 + u_{i\hat{y}}^2 - u_{i\hat{x}}^2 - u_{i\hat{y}}^2) + (K_{pp}/2) \sum_i 4u_{i\hat{x}}^2$$

where $K_{pd}$ and $K_{pp}$ denote the elastic constants. $\hat{x}$ and $\hat{y}$ represent unit vectors in the $x$- and $y$-direction, respectively, $p_{i+\hat{x}/2,\sigma}^\dagger$ and $p_{i+\hat{y}/2,\sigma}^\dagger$ denote the operators for the $p$ electrons at the site $R_i \pm \hat{x}/2$, and in a similar way $p_{i+\hat{y}/2,\sigma}$ and $p_{i+\hat{y}/2,\sigma}$ are defined. $U_d$ denotes the strength of Coulomb interaction between the $d$ electrons. $u_{i\hat{y}}$ and $u_{i\hat{x}}$ represent the variations of the transfer energy $t_{pd}$ and $t_{pp}$, respectively. The number of cells which consist of $d$, $p_x$ and $p_y$ orbitals is denoted as $N$.

The wave function with the inhomogeneous spin structure is made from solutions of the Hartree-Fock Hamiltonian given as

$$H_{trial} = H_{pd}^0 + \sum_{i\sigma} [\delta n_{di} - \sigma (-1)^{\hat{x}i+\hat{y}j} m_i] d_{i\sigma}^\dagger d_{i\sigma},$$

where we have variational parameters $\epsilon_p$ and $\epsilon_d$ in $H_{pd}^0$. In this paper $\delta n_{di}$ and $m_i$ are assumed to have the form

$$\delta n_{di} = - \sum_j \alpha / \cosh \left( (x_i - x_j) \right),$$

where $\alpha$ and $\Delta_{inc}$ are parameters. The inclusion of stripe order parameters considerably improves the ground-state energy. In small clusters the deviation of the energy of stripped state from the exact value is within several percents for the Hubbard model.

The wave function is constructed from the solution of

![FIG. 2: SC condensation energy per site vs the hole density in $t_{pd}$ units, where the parameters are $t_{pp} = 0.4$ and $U_d = 8$. Solid circles and open circles indicate the SC condensation energy for the uniform SC and striped SC, respectively. The lines are fitted by parabola. Squares are obtained for the single-band Hubbard model with the next-nearest transfer $t' = -0.2$ on 12 $\times$ 12 lattice.](image)

![FIG. 3: Lattice structures in the LTT phase (a) and LTO phase (b). The symbol ”+” means that the oxygen atoms move upward and instead ”-” oxygen atoms move downward. ”O” denote the oxygen atom.](image)
Bogoliubov-de Gennes equation given by
\[ \sum_j (H_{ij} u_j^* + F_{ij} v_j^*) = E \lambda u_i, \]

where \( H_{ij} \) and \( F_{ij} \) are \( 3N \times 3N \) matrices including the terms for \( d, p_x \) and \( p_y \) orbitals. The Bogoliubov operators are written in the form
\[ H = \sum_{ij}(u_i^\lambda a_{i\uparrow} + v_i^\lambda a_{i\downarrow}^\dagger) \quad (E^\lambda > 0), \]
\[ H = \sum_{ij}(a_{i\uparrow}^\lambda a_{i\downarrow} + v_i^\lambda a_{i\downarrow}^\dagger) \quad (E^\lambda < 0). \]
\( a_{i\sigma} \) denotes \( d_{i\sigma}, p_{i+\hat{x}/2\sigma} \) or \( p_{i+\hat{y}/2\sigma} \) corresponding to the components of \( u_i^\lambda \) and \( v_i^\lambda \).

Then the wave function is written as \[ \psi = P_G P_N \prod_\lambda \alpha_\lambda \alpha_\lambda^\dagger |0\rangle \]
\[ \propto P_G \{ \sum_{ij} (U^{-1})_{ij} a_{i\uparrow}^\lambda a_{j\downarrow}^\dagger \}^{N_e/2} |0\rangle. \]

Here we show the results in the case without lattice distortions. It has been shown that the striped state is more stable than the uniform SDW state for small hole doping. In Fig.1 the size dependence of the SC condensation energy is shown for the uniform SC in the overdoped region and the striped SC in the underdoped region with the results obtained for the one-band Hubbard model for comparison. The parameters are \( t_{pp} = 0.4 \) and \( U_d = 8 \) in \( t_{pd} \) units. The squares in Fig.1 indicate the SC condensation energy of pure \( d \)-wave state at \( \delta \approx 0.2 \), while the circles are for SC coexisting with stripes at \( \delta = 1/8 \) for \( Q_x = \pi/4 \) evaluated on rectangular lattices \( 8 \times 24, 24 \times 6, 16 \times 8 \) and \( 16 \times 4 \). In both cases the energy obtained through an extrapolation is of the same order as experimental values.

\[ E_{cond} \approx 0.00014 t_{pd} \approx 0.2 \text{meV}, \]
where we have assigned \( t_{pd} \approx 1.5 \text{eV} \). The data in Fig.2 show the SC condensation energy as a function of the hole density. The SC condensation energy per site for the striped SC is reduced as the hole density decreases, while that for pure \( d \)-wave SC remains finite even near half-filling. This suggests that an origin of the decrease of \( T_c \) in the underdoped region lies in the reduction of hole-rich domain where the SC order parameter has finite

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amplitude. Now let us consider the effect of lattice distortion on stripes. In the LTT phase stabilized at low temperatures near 1/8-hole filling, the distortions of the CuO square occur in the manner shown in Fig.3. The LTT phase has a ‘tilting axis’ on which the copper and oxygen atoms never move even in the distorted state. The vertical or horizontal stripes can coexist with the lattice distortions in the LTT phase, being parallel or perpendicular to the tilting axis.

The structural transition from low-temperature orthorhombic (LTO) to LTT phases occurs in LaBaSrCuO and LaNdSrCuO systems. It is not clear a priori what structure is stabilized due to the lattice deformation. We consider the following cases assuming that the stripes are in the $y$-direction:

(A) $u_{ix} = u$, $u_{iy} = 0$, $v_i = 0$,

(B) $u_{ix} = 0$, $u_{iy} = u$, $v_i = 0$,

(C) $u_{ix} = u$, $u_{iy} = u \cos(2Q_i x_i)$, $v_i = u \cos(2Q_i x_i)$,

where $Q_x = 2\pi \delta$ and $u$ is the amplitude of deformation: $u_{ij} = u_{i-j}$, $u = 0$ corresponds to the LTO structure, and the anisotropy in $t_{pd}$ indicates a transition to the LTT phase. $t_{pd}$ increases along the tilt axis compared to the LTO phase. The case (C) corresponds to the structure of mixed LTT-HTT phase. The energy gain per site defined as $\Delta E/N = (E(u = 0) - E(u))/N$ is presented in Fig.4 as a function of $u$ in $t_{pd}$ units. The energy in the case (B) is lower than that in the case (A) indicating that the stripes are parallel to the tilting axis under the rigid LTT structure. The energy in the case (B) is lower than that in the case (A) indicating that the stripes are parallel to the tilting axis under the rigid LTT structure. We simply assume the same elastic energy cost for these types of rotations. The cost of energy due to lattice distortions is assumed to be given by $(K/2)u^2$ for the constant $K$. $K$ is estimated as follows. According to Harrison’s rule, $t_{pd}$ is expected to vary as $d^{-n}$ with $n \approx 7/2$, $d$ being the Cu-O bond length. Since $\delta t_{pd}/t_{pd} = -n\delta d/d$, the elastic energy is estimated as

$$E_{el} = \frac{1}{2}C(2d)^3 \frac{\delta d}{d}^2 = \frac{1}{2}C(2d)^3 \frac{1}{n^2}u^2 \approx \frac{K}{2}u^2. \quad (9)$$

The constant $C$ is estimated as $C \approx 1.7 \times 10^{12}$ dyn/cm$^2 = 1.7$ eV/A$^2$. Since $d \approx 2\AA$, $K$ is of the order of 10eV: $K \approx 8.9$eV. We point out that $E_{el}$ has possibly a linear term in $u$. The $\theta^2$ term, if present in $E_{el}$, is proportional to $u$ since $u \sim 1 - \cos(\theta) \sim \theta^2/2$ where $\theta$ is the tilt angle. The presence of linear term may lead to a first order transition. As shown in Fig.4 the striped state is more stabilized in the LTT phase. We show schematically the stable striped state in the LTT phase in Fig.5 obtained from our VMC evaluations, where the shaded square represents the tilted CuO unit cell rotating around the tilting axis. The LTT-HTT state is more stabilized due to the kinetic energy gain coming from the softening of tilt angles.

In this paper we have investigated the inhomogeneous ground state with the lattice distortions based on the three-band model of high-$T_c$ cuprates using the variational Monte Carlo method. The SC condensation energy decreases as the doping ratio decreases, which is due to the reduction of SC domain in the hole-rich region. The stable striped state has hole-rich arrays being perpendicular to the tilting axis of the lattice distortions in the LTT phase as shown in Fig 5, which can be regarded as the LTT-HTT mixed phase. We thank H. Oyanagi for valuable discussions.

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