A Unique and Simple Approach to Improve Sensitivity in $^{15}$N-NMR Relaxation Measurements for NH$_3^+$ Groups: Application to a Protein-DNA Complex

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Abstract: NMR spectroscopy is a powerful tool for research on protein dynamics. In the past decade, there has been significant progress in the development of NMR methods for studying charged side chains. In particular, NMR methods for lysine side-chain NH$_3^+$ groups have been proven to be powerful for investigating the dynamics of hydrogen bonds or ion pairs that play important roles in biological processes. However, relatively low sensitivity has been a major practical issue in NMR experiments on NH$_3^+$ groups. In this paper, we present a unique and simple approach to improve sensitivity in $^{15}$N relaxation measurements for NH$_3^+$ groups. In this approach, the efficiency of coherence transfers for the desired components are maximized, whereas undesired anti-phase or multi-spin order components are purged through pulse schemes and rapid relaxation. For lysine side-chain NH$_3^+$ groups of a protein-DNA complex, we compared the data obtained with the previous and new pulse sequences under the same conditions and confirmed that the $^{15}$N relaxation parameters were consistent for these datasets. While retaining accuracy in measuring $^{15}$N longitudinal and transverse relaxation measurements, our new pulse sequences for NH$_3^+$ groups allowed an 82% increase in detection sensitivity of $^{15}$N longitudinal and transverse relaxation measurements.

Keywords: dynamics; ion pairs; NH$_3^+$ groups; NMR relaxation; protein side chains

1. Introduction

NMR spectroscopy is one of the most powerful techniques for studying protein dynamics. NMR studies have revealed the functional importance of structural dynamics in many biological molecular processes of proteins (e.g., reviewed in Refs [1–8]). While the vast majority of NMR investigations of protein dynamics have probed motions of either backbone NH or side-chain CH$_3$ groups, NMR investigations on polar or charged side chains remain rare. Recently, there has been significant progress in NMR methods for investigating the dynamics of charged side chains of proteins [9–17]. In particular, NMR methods for Lys side-chain NH$_3^+$ groups have proven to be extremely useful for investigating the dynamics of hydrogen bonding and/or ion pairing [14–25].

Lys side-chain NH$_3^+$ groups of proteins undergo rapid hydrogen exchange with water [26–28]. As a result of this rapid hydrogen exchange, signals from NH$_3^+$ groups in $^1$H-$^{15}$N heteronuclear single-quantum coherence (HSQC) and heteronuclear multiple-quantum coherence (HMQC) spectra are severely broadened [26]. Importantly, this broadening occurs not only in the $^1$H dimension but...
also in the $^{15}\text{N}$ dimension, because rapid hydrogen exchange greatly enhances scalar relaxation of $^{15}\text{N}$ transverse coherence anti-phase with respect to $^1\text{H}$ (e.g., $2N_\text{H}_\text{H}_2$, $4N_\text{H}_\text{H}_2$, and $8N_\text{H}_\text{H}_2$).

To avoid this problem, Iwahara et al. developed $\text{NH}_3^+$-selective heteronuclear in-phase single-quantum coherence (HISQC) and its derivatives [26]. In the HISQC experiment, the in-phase single quantum term $N_x$ or $N_y$ is created at the beginning of the $^{15}\text{N}$ evolution period, and in-phase single-quantum coherence $N_+ (= N_x + iN_y)$ is maintained via the $^1\text{H}$ WALTZ decoupling scheme throughout the evolution period. Evolutions to the anti-phase terms such as $2N_\text{H}_\text{H}_2$, $4N_\text{H}_\text{H}_2$, and $8N_\text{H}_\text{H}_2$ are suppressed to remove the impact of scalar relaxation on line shape of $^{15}\text{N}$ resonances. Scalar relaxation arises from auto-relaxation of the coupled $^1\text{H}$ nuclei [29,30], and substantially increases the relaxation rates of the $2N_\text{H}_\text{H}_2$, $4N_\text{H}_\text{H}_2$, and $8N_\text{H}_\text{H}_2$ terms, compared to the relaxation rates of $N_+$. The scalar relaxation rate $R_{sc}$ for each $^1\text{H}$ nucleus is given by [26]:

$$R_{sc} = \rho_{HH} + k_\text{water}$$  \hspace{1cm} (1)

where $\rho_{HH}$ is the rate for dipole-dipole relaxation with external $^1\text{H}$ nuclei and $k_\text{water}$ is the rate for hydrogen exchange with water. Scalar relaxation rates for the $N_+$, $2N_\text{H}_\text{H}_2$, $4N_\text{H}_\text{H}_2$, and $8N_\text{H}_\text{H}_2$ terms are 0, $R_{sc}$, $2R_{sc}$, and $3R_{sc}$, respectively [31]. Typically, hydrogen exchange is much faster than $\rho_{HH}$ rates and intrinsic $^{15}\text{N}$ relaxation rates for $\text{NH}_3^+$ groups [14–16,26]. Therefore, rapid hydrogen exchange governs relaxation of the anti-phase terms through the scalar relaxation mechanism and severely broadens $^{15}\text{N}$ line shapes of $\text{NH}_3^+$ signals in typical 2D $^1\text{H}-^{15}\text{N}$ correlation spectra. By maintaining in-phase single-quantum terms $N_x$ and $N_y$, and thereby removing the scalar relaxation from the $I_3$ time domain for the $^{15}\text{N}$ dimension, the HISQC experiment drastically improved observation of $^1\text{H}-^{15}\text{N}$ cross peaks from $\text{NH}_3^+$ groups in sensitivity and resolution [26]. Since then, many NMR pulse sequences for $\text{NH}_3^+$ groups have implemented the principle of HISQC, and minimized the adverse impacts of scalar relaxation of anti-phase terms with respect to $^1\text{H}$ nuclei [14–17,26,32].

Nevertheless, relatively low sensitivity due to rapid hydrogen exchange has been a major practical problem in NMR experiments for Lys side-chain $\text{NH}_3^+$ groups of proteins. While some side-chain $\text{NH}_3^+$ groups exhibit relatively slow hydrogen-exchange rates due to hydrogen bonds or ion pairs [26,33], many other $\text{NH}_3^+$ groups exhibit very rapid hydrogen-exchange rates that severely broaden $^1\text{H}$ resonances. Due to this problem, NMR experiments on protein side-chain $\text{NH}_3^+$ groups are often conducted at relatively low pH (typically pH 4.5–6.0) and low temperature (typically, 2–25 °C) to observe a larger number of signals with stronger intensity [32,34]. In these NMR experiments, co-axial NMR tubes that separate lock solvent (usually, D$_2$O) from a sample solution are typically used to avoid isotopically different species (i.e., ND$_2$$^+$, and ND$_2$$^+$, and ND$_3$$^+$) of $\text{NH}_3^+$ groups. The use of co-axial tubes further decreases sensitivity due to a smaller sample volume and multilayer glass walls. Thus, sensitivity improvement would be desirable for NMR experiments on $\text{NH}_3^+$ groups, especially for quantitative experiments such as $^{15}\text{N}$ relaxation measurements.

To address these practical needs, we present a unique and simple approach to improve sensitivity in $^{15}\text{N}$ relaxation measurements on protein side-chain $\text{NH}_3^+$ groups. Our approach involves only minor modifications of the existing pulse sequences. Nevertheless, the pulse sequences implementing this approach significantly improve the detection sensitivity, while maintaining the accuracy in the $^{15}\text{N}$ relaxation measurements on $\text{NH}_3^+$ groups.

2. Results

Figure 1 shows the optimized NMR pulse sequences for measuring $^{15}\text{N}$ $R_1$ and $R_2$ relaxation and heteronuclear NOE of $\text{NH}_3^+$ groups. In the description below, using the product operator formalism [35] for AX$_3$ spin systems, we first explain the previous approach that resolves problems arising from undesired anti-phase or multi-spin-order components of $^{15}\text{N}$ magnetizations of $\text{NH}_3^+$ groups in $^{15}\text{N}$ relaxation measurements. Then, we describe our new approach to improving
sensitivity and eliminating undesired components in $^{15}$N relaxation measurements, showing data that demonstrate the effectiveness of this approach.

![Figure 1](image_url)

**Figure 1.** Pulse sequences for the $^{15}$N relaxation measurement on lysine side-chain NH$_{3}^+$ groups. The key elements in the current work are indicated in red. Thin and bold bars in black represent hard rectangular 90° and 180° pulses, respectively. Water-selective half-Gaussian (2.1 ms) and soft-rectangular (1.2 ms) 90° pulses are represented by half-bell and short-bold shapes, respectively. Unless indicated otherwise, pulse phases are along $x$, and the carrier position for $^1$H was set to the position of the water resonance. The $^{15}$N carrier position was set to 33.1 ppm. A gray bell-shape for $^{15}$N represents an r-SNOB [36] 180° pulse (1.0 ms) selective to Lys side-chain $^{15}$N nuclei. The delays $\tau_d$ and $\tau_b$ were 2.7 ms and 1.3 ms, respectively. Quadrature detection in the $t_1$ domain was achieved using States-TPPI, incrementing the phase $\phi_1$. Pulsed field gradients (PFGs) were optimized to minimize the water signal. (a) $^{15}$N $R_1$ measurement. Although it is not essential owing to negligible CSA-DD cross correlation for NH$_{3}^+$, a $^1$H 180° pulse, which does not affect H$_2$O resonance, was applied every 10 ms during the delay $T_1$ for longitudinal relaxation. Phase cycles: $\phi_1 = (2y, 2y(-y), \phi_2 = (y, -y), \phi_3 = (4x, 4(-x)), \phi_4 = (8y, 8(-y)))$, and receiver = ($x, -x, -x, x, 2(-x, x, x, -x), x, -x, -x, x$); (b) $^{15}$N $R_2$,$\varphi_1$ measurement. The RF strength for $^{15}$N pulses for the CPMG scheme was 5.4 kHz. The $^1$H carrier position was shifted to 7.8 ppm right after the PFG $g_5$ and set back to the position of water resonance right after the PFG $g_5$. The RF strength $\omega_{\text{CPMG}}/2\pi$ of $^1$H CW during the CPMG was set to 4.3 kHz, which was adjusted to satisfy $\omega_{\text{CPMG}}/2\pi = k_{\text{CPMG}}$ (k, integer) [37]. The delays $\xi_1$ and $\xi_2$ are for alignment of $^1$H magnetization and given by $\xi_1 = 1/\omega_{\text{CPMG}} - (4/\pi)\tau_{900}$ and $\xi_2 = \tau_{90N} - (2/\pi)\tau_{90H}$ [37,38], in which $\tau_{90}$ represents a length of a relevant 90° pulse. Phase cycles: $\phi_1 = (4y, 4(-y)), \phi_2 = (8y, 8(-y)), \phi_3 = x, \phi_4 = (x, -x), \phi_5 = (2y, 2(-y)), \phi_6 = (2x, 2(-x)), \phi_7 = (2(-y), 2y))$, and receiver = ($x, -x, -x, 2(-x, x, x, -x), x, -x, -x, x$); (c) Heteronuclear $^1$H,$^{15}$N NOE measurement. Measurement with $^1$H saturation (5 s) was performed with a train of $180^\circ x$ and $180^\circ(-x)$ pulses (RF strength, 11 kHz) at an interval of 10 ms. The $^1$H carrier position was at 7.8 ppm during the $^1$H saturation period. The reference spectrum was measured without the scheme in the bracket. The recycle delay (including the saturation period) was set to 18 s for a 750-MHz spectrometer. Phase cycles: $\phi_1 = (y, -y), \phi_2 = (4x, 4y, 4(-x), 4(-y)), \phi_3 = (2x, 2(-x))$, and receiver = ($x, -x, -x, x, -x, x, -x$, $x$); (d) Efficiency in coherence transfers as a function of the delay $\tau_b$ calculated using Equations (2) and (3) with $^{1H}_{N_{H}}$ $= 74$ Hz and $^1$H 180° pulse length of 20 $\mu$s. The results for the $N_b$ and $4N_b$,$H_2$,$H_2$ terms are shown in solid and dotted lines, respectively. Red and green arrows indicate the values of the delay $\tau_b$ in the current and previous pulse sequences, respectively.

2.1. Previous and Current Approaches to Eliminating the Adverse Effects of Multi-Spin Order Terms

The first step for measuring $^{15}$N longitudinal ($R_1$) and transverse ($R_2$) relaxation rates is to create the $^{15}$N in-phase single-quantum term via coherence transfer from $^1$H to $^{15}$N nuclei through a refocused INEPT scheme [39]. With regard to NH$_{3}^+$ groups, the product operator terms $N_2$, $2N_2$,$H_2$, $4N_2$,$H_2$,$H_2$, and $8N_2$,$H_2$,$H_2$,$H_2$ are generated in the period of $2\tau_b$ in the first refocused INEPT scheme of our pulse sequence for $^{15}$N $R_1$ and $R_2$ measurements (Figure 1a,b). Because the only term of interest among
them is $N_z$, any effects of the other three terms should be eliminated in these relaxation measurements. The $2N_y H_z$ and $8N_y H_z H_z$ terms are eliminated by the pulsed field gradient (PFG) $g_4$ after the $^1H$ $90^\circ(-\chi)$ and $^{15}N$ $90^\circ(y)$ pulses at the end of the refocused INEPT scheme. These $90^\circ$ pulses convert the $N_x$ and $4N_x H_y H_z$ terms into $N_z$ and $4N_z H_y H_y$, both of which survive the PFG $g_4$. The $4N_z H_y H_y$ term survives because a PFG alone cannot destroy homonuclear zero-quantum coherence [40]. To avoid any adverse impact of the $4N_z H_y H_y$ term generated in the refocused INEPT scheme, the previous pulse sequences used a value of the time $\tau_b$ that erases the $4N_y H_z H_z$ term, but retains the $N_z$ term. This is possible because coherence transfer to these terms depends differently on the time $\tau_b$. The coefficients of these transfers are given by [39]:

$$f_{CT}(2N_y H_z \rightarrow N_z) = \cos^2 \theta \sin \theta$$

(2)

$$f_{CT}(2N_y H_z \rightarrow 4N_x H_z H_z) = \left(3 \cos^2 \theta - 1\right) \sin \theta$$

(3)

where $\theta = 2\pi f_{N H \tau_b}$ and $1/\tau_{NH}$ represents the one-bond $^1H$-$^{15}N$ scalar coupling constant. The use of the time $\tau_b$ satisfying $3 \cos^2 \theta - 1 = 0$ thus eliminates the $4N_y H_z H_z$ term, but retains the $N_z$ term [15]. This approach was used for $^{13}C R_1$ and $R_2$ relaxation measurements for protein CH$_3$ groups as well [41,42]. Because $1/\tau_{NH}$ is typically ~74 Hz for lysine side-chain NH$_3^{+}$ groups [26], the condition to suppress the $4N_z H_y H_z$ term was achieved using $\tau_b = 2.1$ ms in the original pulse sequences [15]. This condition was also used in the second refocused INEPT scheme for backward coherence transfer, so that any coherence transfer from $4N_x H_y H_z$ to $2N_y H_z$ does not contribute to the observed signals. A practical problem in using the condition of $f_{CT}(2N_y H_z \rightarrow 4N_y H_z H_z) = 0$ is that it also reduces $f_{CT}(2N_y H_z \rightarrow N_z)$ from its maximum level, and thereby weakens signals in the $^{15}N$ relaxation measurements for NH$_3^{+}$ groups (Figure 1d).

In the current work, we eliminate the adverse effects of the $4N_z H_y H_z$ term in a different manner, and maximize $f_{CT}(2N_y H_z \rightarrow N_z)$ to increase sensitivity in $^{15}N$ relaxation measurements for NH$_3^{+}$ groups. As shown in Figure 1d, the signal arising from the $N_z$ term should be strongest when $\tau_b = 1.3$ ms. Although this condition increases the $4N_x H_z H_z$ term generated through the refocused INEPT scheme, our pulse sequences shown in Figure 1 prevent the undesired $4N_x H_z H_z$ term from becoming observable in the $^1H$ detection period $t_1$. This allows us to use $\tau_b = 1.3$ ms and improve sensitivity without compromising accuracy in $^{15}N$ relaxation measurements.

2.2. Assessment of the Sensitivity-Improved $^{15}N$ R$_1$ Experiment for NH$_3^{+}$ Groups

Our pulse sequence for the $^{15}N$ R$_1$ relaxation measurements on NH$_3^{+}$ groups is shown in Figure 1a. This pulse sequence is the same as that in Esadze et al. [15], except that the time $\tau_b$ is set to 1.3 ms instead of 2.1 ms. The $^1H$ $90^\circ(-\chi)$ and $^{15}N$ $90^\circ(y)$ at the end of the first refocused INEPT convert the $N_x$ and $4N_x H_y H_z$ terms into the $N_z$ and $4N_x H_y H_y$ terms. As mentioned above, both of these terms survive the PFG $g_4$, and are subjected to the period $T_r$ for relaxation measurement. For measuring $^{15}N$ R$_1$ relaxation rates, however, only the $N_z$ term should be retained, and any contribution of the $4N_x H_y H_y$ term should be removed. During the period $T_r$, not only longitudinal relaxation, but also cross-correlation of three $^1H$-$^{15}N$ dipole-dipole (DD) relaxation mechanisms occur for the $N_z$ term. The DD-DD cross-correlation causes partial transitions from $N_z$ to $4N_x H_y H_z$ [43]. The composite of water-selective $^1H$ $90^\circ(-\chi)$ and hard $^1H$ $90^\circ(\chi)$ pulses was originally introduced to prevent this term from becoming detectable while maintaining water $^1H$ magnetization along $z$ [15]. However, if a considerable amount of the $4N_x H_y H_y$ term is present at the beginning of the period $T_r$, the composite pulses at the end can partially convert this term into $4N_z H_y H_z$, which can survive the rest of the pulse sequence and become observable through the second refocused INEPT with $\tau_b = 1.3$ ms. This problem in the $^{15}N$ R$_1$ measurement can readily be resolved by taking advantage of rapid relaxation of the $4N_x H_y H_y$ term. Due to rapid hydrogen exchange with water, scalar relaxation of anti-phase and multi-spin order terms of $^{15}NH_3^{+}$ are far faster than the intrinsic $^{15}N$ R$_1$ and R$_2$ relaxation of NH$_3^{+}$ groups [15,26]. Even under the conditions of pH 5.0 and 2 °C, where hydrogen exchange is relatively slow, the relaxation rates of the $4N_z H_y H_z$ term
were ~20–100-fold faster than the relaxation rates of the $N\_z$ term for the Lys side-chain NH$_3^+$ groups of ubiquitin [15]. The relaxation of the 4$N\_zH\_yH\_y$ term should be even faster because of its transverse nature. Therefore, if the minimum duration of period $T\_r$ in the $^{15}\text{N} R\_1$ relaxation experiment is sufficiently long to let the 4$N\_zH\_yH\_y$ term completely decay, the relaxation rates of the $N\_z$ term (i.e., $^{15}\text{N} R\_1$) can be measured without any adverse contribution from the 4$N\_zH\_yH\_y$ term.

We applied this approach to the Lys side-chain NH$_3^+$ groups of the Antp homeodomain-DNA complex at pH 5.8 and 15 °C. The interfacial Lys side chains K46, K55, K57, and K58 of this protein-DNA complex exhibit well-resolved $^1\text{H}-^{15}\text{N}$ cross peaks in the NH$_3^+$-selective $^1\text{H}-^{15}\text{N}$ HISQC spectra (Figure 2). For these NH$_3^+$ groups, we measured $^{15}\text{N} R\_1$ relaxation rates with the previous and current pulse sequences using the same number of scans and data points. In these $^{15}\text{N} R\_1$ measurements, we recorded 2D $^1\text{H}-^{15}\text{N}$ spectra using $T\_r = 100, 200, 400, 600, 900, 1200, 1600$, and $2100$ ms in an interleaved manner. The minimum duration, $T\_b = 100$ ms, is expected to be long enough to let the 4$N\_zH\_yH\_y$ term completely decay through its rapid relaxation. As predicted in Figure 1d, the signals from NH$_3^+$ groups in the spectra recorded with $T\_b = 1.3$ ms showed significantly stronger intensities than in those recorded with $T\_b = 2.1$ ms. Figure 3a shows the signal intensity of the K46 NH$_3^+$ group as a function of $T\_r$. The sensitivity was found to improve by a factor of 1.82 on average, which was consistent with the ratio of $[f_{CT}(2N\_H\rightarrow N\_z)]^2$ at $T\_b = 1.3$ ms and 2.1 ms. The $^{15}\text{N}$ relaxation rates $R\_1$ were determined through nonlinear least-squares fitting with a single exponential function. Table 1 shows the $^{15}\text{N} R\_1$ relaxation rates measured with the previous and current pulse sequences for the Lys NH$_3^+$ groups in the Antp homeodomain-DNA complex. The $^{15}\text{N} R\_1$ rates from the two experiments were virtually the same, within experimental uncertainties. Not surprisingly, improvement in sensitivity led to higher precision in measured $^{15}\text{N} R\_1$ relaxation rates.

![Figure 2. The $^1\text{H}-^{15}\text{N}$ HISQC spectrum recorded at 15 °C for the NH$_3^+$ groups in the complex of $^{15}\text{N}$-labeled Antp homeodomain and unlabeled 15-bp DNA containing a phosphorodithioate at the K46 interaction site. The resonance assignment is based on that for the unmodified DNA complex and unique chemical shift perturbation upon site-specific dithioation (i.e., sulfur substitutions of two non-bridging oxygen atoms) of the DNA phosphate at the K46 interaction site [44].](image)

| Parameters          | K46             | K55             | K57             | K58             |
|---------------------|-----------------|-----------------|-----------------|-----------------|
| $^{15}\text{N} R\_1$ (s$^{-1}$) | $1.093 \pm 0.013$ | $0.637 \pm 0.005$ | $1.035 \pm 0.004$ | $0.363 \pm 0.002$ |
| $^{15}\text{N} R\_1$ (s$^{-1}$) | $1.081 \pm 0.023$ | $0.617 \pm 0.008$ | $1.037 \pm 0.008$ | $0.364 \pm 0.003$ |
| $^{15}\text{N} R\_2,\text{im}$ (s$^{-1}$) | $2.55 \pm 0.10$  | $1.76 \pm 0.07$  | $2.95 \pm 0.04$  | $1.20 \pm 0.03$  |
| $^{15}\text{N} R\_2,\text{im}$ (s$^{-1}$) | $2.74 \pm 0.20$  | $2.05 \pm 0.12$  | $2.76 \pm 0.06$  | $1.14 \pm 0.06$  |
| Heteronuclear NOE $^b$ | $-2.44 \pm 0.12$ | $-2.83 \pm 0.10$ | $-2.54 \pm 0.05$ | $-2.71 \pm 0.05$ |
| Heteronuclear NOE $^c$ | $-2.53 \pm 0.18$ | $-2.75 \pm 0.13$ | $-2.60 \pm 0.08$ | $-2.65 \pm 0.07$ |

$^a$ The experiments were conducted at 15 °C and the $^1\text{H}$ frequency of 750 MHz. Uncertainties were estimated using the Monte Carlo approach based on the noise standard deviation of the spectra. $^b$ Measured with the current pulse sequences shown in Figure 1. $^c$ Measured with the previous pulse sequences [15].
2.3. Assessment of the Sensitivity-Improved $^{15}\text{N} R_2$ Experiment for NH$_3^+$ Groups

Our new pulse sequence for $^{15}\text{N}$ $R_2$ measurements is shown in Figure 1b. This pulse sequence differs from our previous one in two ways. First, $\tau_b = 1.3$ ms is used instead of $\tau_b = 2.1$ ms. Second, a composite of water-selective $^1\text{H} 90^{\circ}(-x)$ and hard $^1\text{H} 90^{\circ}(x)$ pulses is implemented before the PFG $g_5$. This additional component is important for canceling the effects of the 4$N_b$$_2$H$_2$ $z$ term generated through the refocused INEPT scheme. The pulse sequence uses the CW-CPMG scheme together with H$_2$O alignment pulse trains [37]. During the $^{15}$N CPMG spin-echo periods for $^{15}$N transverse relaxation measurements, a $^1\text{H}$ continuous wave (CW) is applied at the $^1\text{H}$ resonances of NH$_3^+$ groups to maintain the in-phase single-quantum term $N_x$ and prevent the anti-phase terms from being produced. Through the $^1\text{H}$ pulse scheme developed by Hansen et al. [37], water $^1\text{H}$ magnetization is aligned to the axis of $^1\text{H}$ CW in the rotating frame to avoid saturation through $^1\text{H}$ RF inhomogeneity and then is brought back to $+z$. Because this scheme does not align the 4$N_b$$_2$H$_2$$_y$ term, the terms arising from it are largely purged due to the RF inhomogeneity of the $^1\text{H}$ CW. However, their component parallel to the CW axis can remain and become the 4$N_b$$_2$H$_2$$_z$ term through the back-alignment scheme after the period $T_r$. Unlike the period $T_r$ in the $^{15}$N $R_1$ relaxation measurement, the period $T_r$ in the $^{15}$N $R_2$ relaxation measurement should be relatively brief, because only a limited number of hard $^{15}$N $180^{\circ}$ pulses can practically be used during the $^{15}$N CPMG scheme. Therefore, this remaining undesired term cannot be purged completely through relaxation. However, the composite of the water-selective $^1\text{H} 90^{\circ}(-x)$ and hard $^1\text{H} 90^{\circ}(x)$ pulses purges this 4$N_b$$_2$H$_2$$_z$ in the same manner as the 4$N_b$$_2$H$_2$$_z$ term arising from DD-DD cross-correlation during the period $T_r$ is canceled in the $^{15}$N $R_1$ measurement.

For the Lys side-chain NH$_3^+$ groups of the Antp homeodomain in complex with 15-bp DNA, we compared the $^{15}$N $R_2$ relaxation data obtained with the old and new pulse sequences under the same conditions. We recorded nine $^1\text{H}$,$^{15}$N spectra in an interleaved manner using $T_r = 4.8, 14.4, 33.6, 48.0, 76.8, 91.2, and 105.6$ ms. Strictly speaking, $^{15}$N transverse relaxation of NH$_3^+$ groups should occur bi-exponentially due to DD-DD cross-correlation [15,41], but the first 30% decay from the maximum can be treated as a mono-exponential decay, as demonstrated by Esadze et al. [15]. Using mono-exponential fitting, the initial rate constants ($R_{2,ini}$) for this $^{15}$N transverse relaxation were determined from the signal intensity as a function of $T_r$. The results from the data obtained with the previous and current pulse sequences are shown in Figure 3b and Table 1. The $R_{2,ini}$ rates from these two datasets are in good agreement. Due to the use of $\tau_b = 1.3$ ms, the signal intensities in the spectra recorded with the current pulse sequence were significantly higher than those in the spectra recorded with the previous pulse sequence. As expected, the gain in intensity in the $^{15}$N $R_2$ experiment was the

![Figure 3](https://example.com/figure3.png)

**Figure 3.** Comparison of the previous [15] and current pulse sequences for measuring $^{15}$N relaxation of NH$_3^+$ groups. (a,b) $^{15}$N longitudinal (Panel a) and transverse (Panel b) relaxation of the K46 NH$_3^+$ group. The vertical axis represents the signal intensity in the two-dimensional spectra measured as a function of the relaxation period $T_r$. Solid lines represent the best-fit curves obtained through nonlinear least-squares fitting with a mono-exponential function; (c) Slices of the K46 NH$_3^+$ signals along the $^1\text{H}$ dimension from the two-dimensional spectra with and without $^1\text{H}$ saturation for the heteronuclear NOE measurements. In each panel, data obtained with the previous and current pulse sequences are shown in blue and red, respectively.
same as that in the $^{15}$N $R_1$ experiment (i.e., 82% increase on average). This improvement in sensitivity led to significantly higher precision in measured $^{15}$N $R_{2,\text{ini}}$ rates.

2.4. Assessment of the Sensitivity-Improved Heteronuclear NOE Experiment for NH$_3^+$ Groups

Figure 1c shows the pulse sequence for heteronuclear NOE measurements for NH$_3^+$ groups that implements the abovementioned approach. As described by Esadze et al. [15], steady states of the $N_z$ and $4N_zH_zH_z$ terms are created through saturation of $^1$H nuclear magnetization via a train of 180° pulses for heteronuclear NOE measurements on NH$_3^+$ groups. The $4N_zH_zH_z$ steady state occurs due to DD-DD cross-correlation that drives transitions between the $N_z$ and $4N_zH_zH_z$ terms [15]. In the original pulse sequence, $\tau_b = 2.1$ ms was used to avoid any contribution of the $4N_zH_zH_z$ term to the observed signals. However, in the current pulse sequence (Figure 3c), the composite of the water-selective $^1$H 90°($-\chi$) and hard $^1$H 90°($\chi$) pulses convert the $4N_zH_zH_z$ term into $4N_zH_yH_y$ immediately before the $^{15}$N 90° pulse leading to the evolution period $t_1$. As described for the $^{15}$N $R_1$ and $R_2$ experiment, the rest of the pulse sequence does not allow the $4N_zH_yH_y$ term to become observable in the $^1$H detection period $t_2$. Therefore, the use of $\tau_b = 1.3$ ms improves sensitivity without compromising the quality of heteronuclear NOE data, though the gain in sensitivity is relatively small because there is only a single refocused INEPT scheme in this pulse sequence.

We compared the heteronuclear NOE data obtained with the previous and current pulse sequences for the Lys NH$_3^+$ groups of the Antp homeodomain-DNA complex. Figure 3c shows $^1$H slices of the 2D $^1$H-$^{15}$N spectra recorded with and without $^1$H saturation in the heteronuclear NOE experiments. The heteronuclear NOE values from the datasets obtained with the previous and current pulse sequences agreed well, as shown in Table 1. As expected, the spectra recorded with the new pulse sequence exhibited an increase in the intensity of each signal compared with those recorded with the previous pulse sequence under the same conditions. The improvement in the sensitivity was by a factor of 1.35 on average for the heteronuclear NOE measurements.

3. Discussion

As demonstrated above, our new pulse sequences improve sensitivity in $^{15}$N relaxation measurements on protein side-chain NH$_3^+$ groups without compromising accuracy in measuring intrinsic $^{15}$N relaxation parameters. By eliminating contributions from the undesired terms and maintaining the maximum level of coherence transfers of the desired terms, this method increased sensitivity by a factor of 1.82 for the $R_1$ and $R_2$ experiments and by a factor of 1.35 for the heteronuclear NOE experiment. Although our current paper shows data for a protein-DNA complex only, a similar degree of improvement is expected for other systems of different sizes because Equations (2) and (3) are independent of the molecular rotational correlation time. The sensitivity gains for the $^{15}$N $R_1$ and $R_2$ experiments are larger because these experiments include two refocused INEPT schemes, whereas the heteronuclear NOE experiment has one. In fact, the relative magnitudes of the sensitivity gains (i.e., $1.82 \approx 1.35^2$) support this explanation. With the current approach, the time for recording the same quality of data can be significantly reduced compared with the previous $^{15}$N relaxation experiments for NH$_3^+$ groups. The total measurement times for $^{15}$N $R_1$ relaxation, $R_2$ relaxation, and heteronuclear NOE experiments on a 0.8 mM protein-DNA complex (17 kDa) were 18, 20, and 26 h, respectively. Note that signal to noise ratios are proportional to $\sqrt{N_s}$, where $N_s$ is the number of accumulated scans per free induction decay (FID). To get the same data quality using the previous pulse sequences by increasing the number of scans, the total measurement times would approximately be tripled for $^{15}$N $R_1$ and $R_2$ measurements and doubled for the heteronuclear NOE measurement. Because rapid hydrogen exchange of NH$_3^+$ groups weakens their $^1$H signals, the improvement in sensitivity in these relaxation experiments is practically helpful. We hope that this approach will facilitate NMR studies of dynamic processes involving hydrogen bonds and ion pairs and help advance our understanding of protein dynamics and its functional roles.
4. Materials and Methods

The complex of the $^{15}$N-labeled Antp homeodomain and unlabeled 15-bp DNA was prepared as described in our previous papers [21,25,44]. The DNA phosphate group at the K46 interaction site was dithioated in the chemical synthesis, as previously described [14,25]. A 370-µL solution of 0.8 mM complex in a buffer of 20 mM sodium phosphate (pH 5.8) and 20 mM NaCl was sealed in a 5-mm outer tube of a co-axial NMR tube system. To avoid the deuterated species of NH$_3^+$ groups (i.e., ND$_2H^+$, ND$_2$H$^+$, and ND$_3^+$), D$_2$O for the NMR lock signal was sealed separately in an inter insert of the co-axial tube. The NMR experiments were performed at 15 °C with an Avance III spectrometer (Bruker BioSpin, Fällanden, Switzerland) operated at the $^1$H frequency of 750 MHz. A TCI cryogenic probe was used for NMR detection. The $^1$H and $^{15}$N acquisition times were 54 ms and 222 ms, respectively. In each experiment, 16 scans were accumulated per FID, and sub-spectra were recorded in an interleaved manner. The NMR data were processed and analyzed using the NMR-Pipe [45] and NMR-View [46] programs. Other experimental details are given in figure captions. The pulse programs and parameter sets for Bruker NMR spectrometers are available upon request via https://scsb.utmb.edu/labgroups/iwahara/software.

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Author Contributions: D.N. prepared the sample of the protein-DNA complex for the NMR experiments; G.L.R.L. and D.E.V. synthesized the DNA strand containing a phosphorodithioate; and J.I. designed the research, conducted the NMR experiments, analyzed the data, and wrote the manuscript.

Conflicts of Interest: The authors declare no conflict of interest.

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