Numerical Simulation of Hydraulic Jump over Rough Beds

Saman Nikmehr1*, Younes Aminpour2

1 Department of Water Science and Engineering, University of Kurdistan, Kurdistan Province, Pasdaran Blvd. Sanandaj, P.O.B. 66177-15175, Iran
2 Department of Hydraulic Structures, Water Research Institute, Ministry of Energy, Tehran, Iran

* Corresponding author, e-mail: s.nikmehr@uok.ac.ir

Received: 15 November 2019, Accepted: 24 February 2020, Published online: 16 March 2020

Abstract

In this study, the hydraulic jumps over rough beds are numerically simulated. In order to calibrate the numerical model, the experimental data were used, which performed in a rectangular flume in various roughness arrangements and different Froude numbers. The effect of the distance (s) and the height (t) of the roughness on different characteristics of the hydraulic jump, including the sequent depth ratio, water surface profile, jump's length, roller's length, and velocity distribution were evaluated and compared. The results showed that the numerical model is fairly well able to simulate the hydraulic jump characteristics. The results also showed that the height and distance of roughness slightly reduced the sequent depth ratios for all Froude numbers. Also, the hydraulic jump length is reduced at the presence of the rough bed. Velocity profiles in different experiments were similar and there was a good agreement between simulated and measured results. Also, increasing the distance and the height of the roughness will slow down the velocity near the bed, increase the shear stress, and increase the gradient of the velocity distribution near the bed.

Keywords

hydraulic jump, numerical simulation, rough beds, FLOW 3D Model

1 Introduction

Hydraulic jump is defined as the sudden transition from a supercritical flow to subcritical condition in a short distance. Furthermore, the hydraulic jump is a phenomenon where the water surface moves upwards through critical depth as kinetic energy is converted to potential energy [1, 2]. Hydraulic jumps are usually used to dissipate excessive energy downstream of the hydraulic structures. Hydraulic jumps on smooth beds have been extensively investigated [3–11]. In these beds, considering the depth and initial flow velocity, and using existing mathematical and empirical relationships, the sequent depth ratio and length of hydraulic jump can be determined.

Until now, many experimental studies have been carried out on the hydraulic jump over a rough bed [12–15]. Rajaratnam [12] is one of the pioneers who widely studied the characteristics of the hydraulic jump on rough bed channels. He introduced the parameter \( k = k_e/y_1 \) representing the effect of bed roughness on the lengths of rollers zone (\( L_r \)) and length of hydraulic jump (\( L_j \)), in which \( k_e \) is the equivalent height of bed roughness and \( y_1 \) is the depth of incoming supercritical flow. Hughes and Flack [13] showed that the boundary roughness decreases the sequent depth and length of hydraulic jump depending on Froude number of incoming flow and relative height of bed roughness. Mohamed [16] used square cross strips and observed 17.4 to 67.4 % decrease on the lengths of hydraulic jump when the Froude number of incoming flow was varied from 4 to 10. Ead et al. [17] reported that the sequent depth of hydraulic jump on corrugated beds is less than that on smooth beds where the lengths of hydraulic jump decreases with 50 % and shear stress increases by 10 times. Tokyay [18] investigated the effect of channel bed corrugations on a hydraulic jump experimentally. In the experiments 2 values of wave steepness, 0.20 and 0.26, were used and the range of Froude numbers was from 5 to 12. The results showed that many factors, such as initial flow depth, supercritical Froude number, height, and wave length of corrugations were effective on the characteristics of the hydraulic jump. Izadjoo and Shafai Bejestan [19] showed that the sequent depths of the hydraulic jump on corrugated beds is 20 % less than that on smooth beds. They reported that the length of hydraulic jump on trapezoidal corrugated beds may decrease by 50 % which is mostly the consequences of roughness crest than the openings. They also concluded
that the shear stress increases by 10 times compared to classic hydraulic jump. Ead and Elsebaie [20] used three sinusoidal, trapezoidal and triangular corrugated rough beds and found that the shape of rough elements does not play an important role on $y_2/y_1$. Carollo and Ferro [21] studied the characteristics of hydraulic jump on a rip-raped bed and concluded that relative sequent depth $y_2/y_1$ decreases with the size of bed roughness. Hassanpour et al. [22] investigated the effect of geometric and hydraulic parameters on energy dissipation and location of the hydraulic jump, in different roughness heights and the divergence of walls. Their results showed that the tail-water depth, the roller length, and the hydraulic jump length on a gradual expansion basin with the rough bed were significantly smaller than that of the classical hydraulic jumps in a rectangular basin with smooth and rough bed. It should be mentioned that Pagliara and Palermo [23] and Felder and Chanson [24] took into consideration the effect of air concentration, bed roughness and channel slope on both sequent depth ratio and other lengths of HJs occurring on rough beds.

A numerical simulation of jump on smooth beds in the rectangular channels has been studied by many researchers. Gharangik and Chaudhry [25] investigated a hydraulic jump by a numerical model. They applied the Boussinesq equations to simulate both the sub and supercritical flows and a hydraulic jump in a rectangular channel having a small bed slope. The numerical study of hydraulic jump on a smooth bed has also been carried out by Zhao et al. [26]. The governing equations are the continuity and momentum for incompressible flow and based on the 2-dimensional $k-\varepsilon$ turbulence model. They used the volume of fluid (VOF) and the scale turbulence model to predict water surface location and horizontal velocity. Sarker and Rhodes [27] studied hydraulic jump on a smooth bed by physical and numerical methods. They used the Re-Normalization Group (RNG) and $k-\varepsilon$ turbulence models in combination with the VOF method for free surface modeling. There was a good agreement between the 2-dimensional computational fluid dynamics (CFD) solution and the physical measurements. Gonzalez and Bombardelli [28] simulated a hydraulic jump on a smooth bed using $k-\varepsilon$ and Large Eddy Simulation (LES) turbulence models. The simulated results were compared with observations of mean flow and turbulence in hydraulic jumps by Liu et al. [29]. Velioglu et al. [30] studied the effects of strip roughness elements on the characteristics of hydraulic jumps such as conjugate depth ratio, jump length and energy dissipation, experimentally and numerically using Flow 3D model. They found that the strip roughness elements have significant effects on the hydraulic jump characteristics. The tail water depth reduction compared to classical jump is 18–20 %. The length of the hydraulic jump is reduced about by 20–25 %. This type of roughness elements induces 2–3 % more energy dissipation than that of a classical hydraulic jump. Abbaspour et al. [31] performed numerical simulation of hydraulic jump on a corrugated bed using both standard $k-\varepsilon$ and RNG models. They used the VOF method for free surface. Their results showed that $k-\varepsilon$ turbulent model and VOF method were suitable for predicting the water surface in the jump on a corrugated bed and the relative error of predicted water surface profiles and measured value were in the range of 1 %–8.6 %. Bayon et al. [32] evaluated the performance of the FLOW-3D and Open FOAM models for simulation of the hydraulic jump in low Reynolds numbers. Comparison of numerical and experimental results showed a good agreement for sequent depth ratio and length of the hydraulic jump for both models. Recently, Viti et al. [33] have presented a literature review on the recent results and future outlook of the numerical simulation of hydraulic jumps.

In addition, the effects of the air concentration on the hydraulic jump characteristics have been studied. Pagliara and Palermo [23] showed that in a two-phase flow (water and entrained air), locating the water surface is problematic and an equivalent flow depth ($d_e$) is commonly used. The depth $d_e$ represents the normal distance from the channel invert reference (effective top ET) to an elevation where the air concentration reaches 90 %. Similarly, Felder and Chanson [24] noticed that the upstream flow was rough and aerated for several flow conditions affecting the accuracy of the flow depth recording, and that the upstream depth, $d_u$, defined in terms of the equivalent clear-water depth, $d$, determined based upon void fraction measurements.

Many previous studies have focused on the hydraulic jumps over smooth beds and few studies have investigated the effect of rough bed on the characteristics of the hydraulic jump numerically. Therefore, in this paper, important the hydraulic jump characteristics such as the water surface profile, jump and roller length, and velocity distribution were investigated based on a detailed comparison with an experimental study. Furthermore, this paper will answer some important questions about the ability of Flow-3D numerical model to simulate the hydraulic jump characteristics and how much the differences between the experimental data and numerical outputs.
2 Material and methods

2.1 Experimental data

The experimental were conducted at the Water Research Center, Department of Irrigation and Reclamation Engineering (University of Tehran) in a metal-glass flume of rectangular cross section. The flume has the dimensions of 0.25 m width, 0.5 m depth, and 12 m length with a slope of 0.002. The flume is equipped with a sluice gate at the entrance and discharge was measured by a triangular weir placed at the end of it.

The initial and secondary depths of the hydraulic jump were measured using a point gauge with a precision of 0.1 mm. Longitudinal water surface profile of the hydraulic jump was determined by means of photographing and digitizing using Grapher software (ver. 7.0.1870, Golden Software, Colorado 80401). In order to stabilize the hydraulic jump and form the initial depth, a crack valve was used at a distance of 17 cm downstream of the sliding valve, to adjust the secondary depth of the hydraulic jump. Fig. 1 illustrates the flow conditions and hydraulic jump characteristics. In addition, the effective parameters and the range of variations are presented in Fig. 2 and Table 1.

2.2 Numerical modelling

In the present research, hydraulic jump on a corrugated bed was numerically studied for different Froude numbers using the 2-phase flow theory. The water surface location was determined with the VOF method. RNG turbulent models were used to simulate the hydraulic jump on corrugated bed. Flow-3D is a CFD model which uses VOF method. Flow 3D (Ver. 11.2.0.16, Flow Science) model gives satisfactory results on free surface flows by solving the Reynolds-averaged Navier-Stokes (RANS) equations over the computational domain [34]. For three-dimensional simulation of the flow in an incompressible state, this model solves the Reynolds-averaged Navier-Stokes (RANS) equations using a finite volume method on a grid domain. These equations in Cartesian coordinate system are as follows:

\[
\begin{align*}
\frac{\partial p}{\partial t} + \frac{\partial(puA_x)}{\partial x} + \frac{\partial(pvA_y)}{\partial y} + \frac{\partial(pwA_z)}{\partial z} &= R_{SOR} + R_{TDF}, \\
\frac{\partial u}{\partial t} + \frac{1}{V_f} (uA_x \frac{\partial u}{\partial x} + vA_y \frac{\partial u}{\partial y} + wA_z \frac{\partial u}{\partial z}) &= \frac{1}{\rho} \frac{\partial P}{\partial x} + G_x + f_x, \\
\frac{\partial v}{\partial t} + \frac{1}{V_f} (uA_x \frac{\partial v}{\partial x} + vA_y \frac{\partial v}{\partial y} + wA_z \frac{\partial v}{\partial z}) &= \frac{1}{\rho} \frac{\partial P}{\partial y} + G_y + f_y, \\
\frac{\partial w}{\partial t} + \frac{1}{V_f} (uA_x \frac{\partial w}{\partial x} + vA_y \frac{\partial w}{\partial y} + wA_z \frac{\partial w}{\partial z}) &= \frac{1}{\rho} \frac{\partial P}{\partial z} + G_z + f_z,
\end{align*}
\]

where, \(u, v, w\) are velocity components; \(A_x, A_y,\) and \(A_z\) are the flow areas, \(G_x, G_y,\) and \(G_z\) are gravitational accelerations; \(f_x, f_y,\) and \(f_z\) are accelerations due to the viscosity; \(\rho\) is the fluid density; \(R_{SOR}\) is the mass source; \(R_{TDF}\) is the source of turbulent diffusion; \(V_f\) is the flow volume and \(P\) is the pressure. In all the simulations, the six-sided boundary

![Fig. 1 (a) The view of the laboratory flume used for the experiments; (b) A view of the used roughness; and (c) and (d) A schematic view of the hydraulic jump formed downstream of the sluice gate [35]](image1)

![Fig. 2 The schematic view of formed hydraulic jump over the rough bed and its characteristics](image2)

| Table 1 Studied Parameters and range of changes |
|---------------------------------------------|
| Effective Parameters | Symbol | Range of Variations |
|-----------------------|--------|---------------------|
| Initial depth of the hydraulic jump | \(y_1\) (m) | 0.012–0.016 |
| Sequent depth of the hydraulic jump | \(y_2\) (m) | 0.094–0.222 |
| Hydraulic jump length | \(L_j\) (m) | 0.320–0.800 |
| Roller region length | \(L_r\) (m) | 0.150–0.460 |
| Flow Discharge | \(Q\) (m³/s) | 0.006–0.015 |
| Initial Froude Number | \(F_{Fr}\) (-) | 5.010–13.700 |
| Roughness Height | \(t\) (m) | 0.020, 0.030, 0.040 |
| The distance between two roughnesses | \(s\) (m) | 0.040, 0.060, 0.080 |
conditions were introduced into the Flow 3D model: $X_{\text{min}}$ (upstream sluice face, the upstream depth was applied to the model at a distance of 1.2 m), $X_{\text{max}}$ (end of laboratory flume, tail-water depth was introduced as a downstream boundary condition). $Y_{\text{min}}$ and $Y_{\text{max}}$ (The left and right walls were introduced to the model as wall boundary condition). For $Z_{\text{min}}$ and $Z_{\text{max}}$, the wall and symmetry boundary conditions were used, respectively. After defining the experimental conditions to the numerical model, the model was run. Fig. 3 shows the output numerical simulation for different flow conditions.

2.3 Verification using experimental data

After obtaining the outputs of the numerical model for the water surface profile, the numerical results were compared with the experimental data of Ravar [35]. In this study, the water surface profile (initial and secondary depths of the hydraulic jump) was used for verification of the numerical model. It should be noted that the water surface profile was determined using the VOF method. Numerical simulations were calibrated with both discharge and water surface profiles. Fig. 4 shows a comparison between the experimental and numerical results of two parameters of discharge and water surface profiles. Different characteristics of hydraulic jump including the water surface profile, sequent depth ratio, the length of hydraulic jump, and velocity distributions over the rough bed with different arrangement and the results are summarized in the following (Table 2).

3 Results

3.1 Water surface profile

By analyzing the results of the model output, the water surface profiles were obtained for different tests are shown in Fig. 5. By plotting non-dimensional values of $(x/L_j)$ vs. $(y-y_j)/(y_2-y_1)$, the water surface profiles were obtained for different experiment (Fig. 6). According to this figure, the water surface profiles are similar in different conditions. In addition, the following relationship was obtained for dimensionless water surface profiles:

$$\frac{(y-y_j)}{(y_2-y_1)} = \frac{x/L_j}{0.026 + \left(\frac{x/L_j}{2.143}\right)^2},$$

where, $x$ and $y$ are the horizontal and vertical distances, $y_j$ and $y_2$ are the initial and secondary depths of the hydraulic jump and $L_j$ is the hydraulic jump’s length. This relationship predict the dimensionless water surface profile with high accuracy (MAE = 6.02 %).
3.2 Sequent depth ratio

Fig. 7(a) shows the values of $y_2/y_1$ against $Fr_1$. A comparison between numerical and experimental measurements of the sequent depth ratio over the rough bed and the data obtained from the Belanger equation (Belanger [36]; Chanson [37]) is shown in Fig. 6(b). As can be seen, the simulation and theoretical results are compatible with each other, and the existing differences can be applied to the errors in the numerical solution, and the hypotheses refer to the Belanger's equation. The mean error between the numerical and experimental simulation results is about 4.8%. Existing differences can be attributed to the errors in the numerical simulation and the hypotheses considered in the Belanger's equation.

In general, the results showed that the height and distance of roughness reduce the sequent depth ratio for all Froude numbers and this trend is intensified by increasing the Froude number. In other words, reduction of the boundary layer thickness will further increase the effect of

| Run No. | $t$ (m) | $s$ (m) | $Q$ (m$^3$/s) | Experimental | Simulated |
|---------|---------|---------|---------------|--------------|-----------|
|          | Fr$_1$ | $y_1$ (m) | Fr$_1$ | $L_j$ (m) | Fr$_1$ | $y_1$ (m) | Fr$_1$ | $L_j$ (m) |
| 1       | 0.008  | 5.01   | 0.10 | 0.17 | 39.6  | 0.016 | 0.10 | 0.13 | 35.0 |
| 2       | 0.011  | 7.04   | 0.15 | 0.28 | 0.55  | 0.016 | 0.15 | 0.23 | 0.50 |
| 3       | 0.015  | 9.86   | 0.21 | 0.37 | 0.79  | 0.016 | 0.20 | 0.34 | 0.66 |
| 4       | 0.006  | 6.03   | 0.09 | 0.15 | 0.38  | 0.012 | 0.10 | 0.18 | 0.32 |
| 5       | 0.014  | 13.70  | 0.22 | 0.46 | 0.80  | 0.012 | 0.21 | 0.42 | 0.76 |
| 6       | 0.008  | 5.09   | 0.10 | 0.16 | 0.32  | 0.016 | 0.11 | 0.13 | 0.27 |
| 7       | 0.011  | 7.16   | 0.15 | 0.23 | 0.54  | 0.016 | 0.16 | 0.20 | 0.47 |
| 8       | 0.013  | 8.89   | 0.19 | 0.31 | 0.67  | 0.016 | 0.18 | 0.28 | 0.61 |
| 9       | 0.009  | 6.04   | 0.12 | 0.17 | 0.33  | 0.016 | 0.13 | 0.17 | 0.31 |
| 10      | 0.011  | 7.10   | 0.15 | 0.20 | 0.36  | 0.016 | 0.16 | 0.21 | 0.31 |
| 11      | 0.012  | 8.13   | 0.17 | 0.22 | 0.43  | 0.016 | 0.16 | 0.23 | 0.40 |
| 12      | 0.009  | 9.09   | 0.15 | 0.23 | 0.54  | 0.012 | 0.14 | 0.23 | 0.49 |
| 13      | 0.012  | 11.87  | 0.19 | 0.28 | 0.62  | 0.012 | 0.18 | 0.26 | 0.56 |

Fig. 5 The water surface profiles obtained for different tests

Fig. 6 Dimensionless water surface profiles over the rough bed for all simulations

3.2 Sequent depth ratio

Fig. 7(a) shows the values of $y_2/y_1$ against $Fr_1$. A comparison between numerical and experimental measurements of the sequent depth ratio over the rough bed and the data obtained from the Belanger equation (Belanger [36]; Chanson [37]) is shown in Fig. 6(b). As can be seen, the simulation and theoretical results are compatible with each other, and the existing differences can be applied to the errors in the numerical solution, and the hypotheses refer to the Belanger's equation. The mean error between the numerical and experimental simulation results is about 4.8%. Existing differences can be attributed to the errors in the numerical simulation and the hypotheses considered in the Belanger's equation.

In general, the results showed that the height and distance of roughness reduce the sequent depth ratio for all Froude numbers and this trend is intensified by increasing the Froude number. In other words, reduction of the boundary layer thickness will further increase the effect of

Table 2 List of simulations performed in this study

| Run No. | $t$ (m) | $s$ (m) | $Q$ (m$^3$/s) | Fr$_1$ | $y_1$ (m) | Fr$_1$ | $L_j$ (m) |
|---------|---------|---------|---------------|-------|-----------|-------|-----------|
|          |         |         |               | Fr$_1$ | $y_1$ (m) | Fr$_1$ | $L_j$ (m) |
| 1       | 0.008  | 5.01   | 0.10 | 0.17 | 39.6  | 0.016 | 0.10 | 0.13 | 35.0 |
| 2       | 0.011  | 7.04   | 0.15 | 0.28 | 0.55  | 0.016 | 0.15 | 0.23 | 0.50 |
| 3       | 0.015  | 9.86   | 0.21 | 0.37 | 0.79  | 0.016 | 0.20 | 0.34 | 0.66 |
| 4       | 0.006  | 6.03   | 0.09 | 0.15 | 0.38  | 0.012 | 0.10 | 0.18 | 0.32 |
| 5       | 0.014  | 13.70  | 0.22 | 0.46 | 0.80  | 0.012 | 0.21 | 0.42 | 0.76 |
| 6       | 0.008  | 5.09   | 0.10 | 0.16 | 0.32  | 0.016 | 0.11 | 0.13 | 0.27 |
| 7       | 0.011  | 7.16   | 0.15 | 0.23 | 0.54  | 0.016 | 0.14 | 0.20 | 0.47 |
| 8       | 0.013  | 8.89   | 0.19 | 0.31 | 0.67  | 0.016 | 0.18 | 0.28 | 0.61 |
| 9       | 0.009  | 6.04   | 0.12 | 0.17 | 0.33  | 0.016 | 0.13 | 0.17 | 0.31 |
| 10      | 0.011  | 7.10   | 0.15 | 0.20 | 0.36  | 0.016 | 0.16 | 0.21 | 0.31 |
| 11      | 0.012  | 8.13   | 0.17 | 0.22 | 0.43  | 0.016 | 0.16 | 0.23 | 0.40 |
| 12      | 0.009  | 9.09   | 0.15 | 0.23 | 0.54  | 0.012 | 0.14 | 0.23 | 0.49 |
| 13      | 0.012  | 11.87  | 0.19 | 0.28 | 0.62  | 0.012 | 0.18 | 0.26 | 0.56 |
roughness and intensify reduction of the secondary depth ratio. This slight decrease can also be attributed to the increase in the vortex moving between the roughnesses, which increases with the increasing the Froude number.

3.3 The hydraulic jump length

Hydraulic jump length is defined as the distance between the two cross-sections with the conjugate depths. In the experimental study, the length of the hydraulic jump is equal to the distance between the beginning of the jump and a point where the water surface fluctuations fades. Similarly, in the results of numerical simulations, the same method was used to determine the length of the hydraulic jump. Fig. 8(a) shows the experimental data and numerical results for relative length of the hydraulic jump ($L_j/y_1$) against the Froude numbers ($Fr_1$) on the rough bed. In addition, Fig. 8(b) shows a comparison between the experimental data and numerical results for relative length of the hydraulic jump. According to these figures, the changes in the experimental and numerical results are in a good agreement (MARE = 10.85 %). It should be noted that numerical model have underestimated the hydraulic jump length compared to experimental conditions.

In addition, with increasing the height and distance between roughnesses, the length of the hydraulic jump decreased significantly. Increase of the Froude number may intensify this trend, so that for the maximum Froude numbers the length of the hydraulic jump decreased about 28 %. This reduction is due to the formation of vortex flow, and its effects may appear as a reduction in the pressure and an increase in the mixing the air bubbles.

3.4 Roller length

Fig. 9(a) shows the relative roller length of the hydraulic jump ($L_r/y_1$) against $Fr_1$ in different the heights and distances between roughnesses. According to Fig. 9(b), it can be seen that the numerical model is well able to predict the relative roller length with MARE = 10.59 %. In this study, the roller length was measured similar to the length of the hydraulic jump. The results showed that the reduction in the roller length depends mainly on the space of the roughnesses, in which the vortex flow forms. For this reason, with increasing the height and distance between the roughnesses, the roller length decreases by about 36 %. In general, as mentioned by Pagliara and Palermo [23] and Felder and Chanson [24], the hydraulic jump characteristics are
deeply affected by the air concentration. This is a fundamental issue, having a significant impact on the generalizability and correctness of the results. This can also be considered as an important reason for observing errors between experimental and simulated results.

3.5 Velocity distribution

In the experimental study, the velocity distributions have not been measured. However, according to the calibration of the numerical model based on experimental data (water surface profiles), one can rely on other outputs of numerical model. Typical velocity distribution in hydraulic jump is shown in Fig. 10. In this figure, the parameter $b$ is the value of $y$ at which $u = 0.5 u_{\text{max}}$, and $\delta$ represents the position of maximum velocity.

In this study, the velocity distribution was simulated at 4 sections during the hydraulic jump. These sections are different for each arrangement of roughness (due to the difference in the length of the hydraulic jump in each mode). Furthermore, these sections were measured on top of the roughness elements. The velocity distribution in hydraulic jumps at different distances from the beginning of the hydraulic jump and different Froude numbers are shown in Figs. 11 to 13. In these figures, $x$ represents the distance from the beginning of the hydraulic jump. Since the shear stress decreases with increasing distance from the beginning of the hydraulic jump and the flow depth increases, the maximum velocity at each section will decrease along the hydraulic jump.

The velocity distribution in the hydraulic jumps is usually shown by plotting the ratio of $y/b$ versus $u/u_{\text{max}}$. As can be seen in Fig. 14, the velocity distributions in hydraulic jumps is the same for all the Froude numbers. In addition, increasing the distance and the height of the roughness will reduce the velocities near the bed and increase the gradient of the velocity distribution near the bed. In other words, on the rough beds, the maximum velocity will occur in a higher distance from the bed than that of the smooth beds, indicating the increase in the thickness of the inner layer of the horizontal velocity distribution.

Fig. 15 presents the dimensionless velocity distribution in hydraulic jump for all the rough beds. As can be seen, the velocity distributions follow a similar trend for all conditions, and the change in height and distance does not significantly affect the velocity distributions.

The ratio $b/y_1$ are plotted against the ratio of $x/y_1$ (Fig. 16). The parameter $b$ represents the value of $y$ at which $u = 0.5 u_{\text{max}}$. As can be seen, by increasing the distance from the beginning of the hydraulic jump, the parameter $b$ increases. This shows that with increasing distance from the beginning of the hydraulic jump, the position of maximum velocity (i.e., the thickness of the inner layer of the horizontal velocity distribution) increases. It should be noted that the height and distance of roughness have no significant effect on this parameter.

The following equation was also developed to estimate the relative position of maximum velocity ($b/y_1$) as a function of ($x/y_1$):

$$\frac{b}{y_1} = \text{function of } \frac{x}{y_1}$$
in which, \( x \) is the desired distance from the beginning of the hydraulic jump.

4 Conclusions

In this study, the characteristics of hydraulic jump including water surface profile, sequent depth ratio, hydraulic jump length, roller length and velocity distribution over rough beds with trapezoid blocks were numerically investigated using Flow-3D model, which it was not completely considered in the previous studies. In order to simulate the flow turbulence and hydraulic jump formed over the trapezoid blocks, a RNG method. The results of this study are briefly summarized as follows:

- The water surface profiles of the hydraulic jump are similar for different tests. This result was also obtained in the experimental study.
- The results showed that the value of MARE between the simulated and measured sequent depth ratios are 4.8 %. The results also showed that the height and distance of roughness slightly reduced the sequent depth ratios for all Froude numbers.
- The value of MARE between the experimental and the numerical results for the hydraulic jump length was about 10.85 %. This difference can be due to the difference in the criterion used for measuring the length of the hydraulic jump. In addition, with increasing the height and distance between roughnesses, the lengths of the hydraulic jump decreased significantly.
Fig. 12 Velocity distribution at different distances from the beginning of the hydraulic jump for different Froude numbers and different flow discharges ($t = 0.04$ m, $s = 0.08$ m)

Fig. 13 Velocity distribution at different distances from the beginning of the hydraulic jump for different Froude numbers and different flow discharges ($t = 0.03$ m, $s = 0.06$ m)
By analyzing the velocity data during the hydraulic jump, it was found that the maximum velocity at each section will decrease along the hydraulic jump. Non-dimensional velocity distribution in the hydraulic jumps is the same for all Froude numbers. In addition, increasing the distance and the height of the roughness will slow down the velocity near the bed and increase the gradient of the velocity distribution near the bed.

Fig. 14 Comparison of the dimensionless velocity distribution in rough beds

(a) $t = 0.02 \text{ m}, s = 0.04 \text{ m}$

(b) $t = 0.03 \text{ m}, s = 0.06 \text{ m}$

(c) $t = 0.04 \text{ m}, s = 0.08 \text{ m}$

Fig. 15 Dimensionless velocity distribution in the hydraulic jump for different Froude numbers 280 in rough beds

Fig. 16 The ratio $b/y_1$ vs. $x/y_1$ for different rough beds

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