Three-dimensional numerical simulation of stepped dropshaft with different step shapes
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ABSTRACT

The deep tunnel system is increasingly used worldwide for stormwater conveyance and storage, providing a robust and effective means of preventing urban waterlogging. In the system, the dropshaft, with the function of conveying stormwater to the deep tunnels underground, often runs under conditions of high falling head and large discharge. Based on the standard stepped dropshaft, a blade-shaped stepped dropshaft was proposed in order to control the potential standing wave and improve discharge capacity. Its hydraulic characteristics in respect of flow pattern, flow rate distribution, time-averaged pressure and energy dissipation were investigated by numerical simulation. Compared with the standard stepped dropshaft, the blade-shaped stepped dropshaft generated a more uniform flow rate distribution in the radial direction, therefore effectively decreasing the height of the standing wave near the external wall. The negative pressure areas that easily existed on the vertical wall of steps were well controlled. The energy dissipation of the blade-shaped stepped dropshaft was as high as that of the standard stepped dropshaft. Therefore, the blade-shaped stepped dropshaft could be a preferable design for the deep tunnel system.

Key words | blade-shaped, hydraulic characteristic, numerical simulation, stepped dropshaft

HIGHLIGHTS

- Report a new kind of stepped dropshaft in deep tunnel systems, namely the blade-shaped stepped dropshaft.
- Obtain the hydraulic characteristics of the blade-shaped stepped dropshaft in respect of flow pattern, flow rate distribution, time-averaged pressure and energy dissipation.
- Compare with the standard stepped dropshaft to obtain the influence of the step shape on the hydraulic characteristics.

INTRODUCTION

With the rapid development of urbanization, the increase of population and the decrease of water area have led to great pressure on the sewer and storage tunnels of the city (Vasconcelos & Wright 2006). Due to global climate change, the pressure may further increase with frequent intense rain events leading to urban water problems such as waterlogging (Zhang et al. 2016). Deep storage tunnel systems, generally including several dropshafts and underground tunnel, can be effective in preventing the urban water problems, and have been implemented or is under consideration in many cities (Guo & Song 1990; Changnon 2010). In this system, the stormwater runoff can be diverted to the deep tunnels through dropshafts, and then diverted to the river or sewage plant, or temporarily stored in the
tunnels (Guo & Song 1991). The dropshaft often runs with the conditions of high falling head and large discharge, which may lead to operational issues related to energy dissipation and cavitation, and even cause disasters like geysers (Rajaratnam et al. 1997; Vasconcelos & Wright 2011). Therefore, it is important to take into account the hydraulic characteristics in the dropshaft.

Dropshafts can be divided into four typical types: the vortex dropshaft, baffle dropshaft, plunging dropshaft and helicoidal-ramp dropshaft (He et al. 2017). Previous research indicated that the helicoidal-ramp dropshaft shows a better performance for high falling head and large discharge with good energy dissipation and exhaust effect (Kawasaki et al. 1999). It includes a vertical circular dropshaft and a continuous helicoidal ramp attached to the dropshaft wall. The water can skim down along the ramp, and the air can be released from the center hollow column without any extra ventilation device (Ansar & Jain 1997). The energy dissipation occurs as distributed friction loss along the helicoidal ramp without any extra energy dissipators (Kennedy et al. 1988).

Inspired by the helicoidal-ramp dropshaft, Wu et al. (2018) proposed a stepped dropshaft in order to provide well aeration to decrease the risk of cavitation erosion and to further increase the energy dissipation. Figure 1(a) is a definition sketch of the standard stepped dropshaft. By changing the ramps to successive steps, the flow skimming over the ramps could be converted to successive nappes impacting steps with vortices beneath them. On the one hand, the energy dissipation can increase due to nappe breakup in the air and impact on the steps as well as full or partial hydraulic jumps. On the other hand, the intense turbulence resulting from the vortices increases the self-aeration from the free surface, which can decrease the risk of cavitation.

Some investigations have been focused on the hydraulic characteristics of the stepped dropshaft. Wu et al. (2018) indicated that there are two typical flow regimes in the stepped dropshaft: nappe flow and skimming flow, which are similar to that in the stepped spillway (Qian & Wu 2020; Wu et al. 2020). Compared with the stepped spillway, the flow affected under the centrifugal force is more complicated for the stepped dropshaft. Wu et al. (2017) showed that the height of the standing wave near the external wall decreases with the increase in the approach flow Froude number, but increases with the increase in the curvature of the dropshaft. The influence of the geometric parameter of the stepped dropshaft on the hydraulic characteristics has been investigated by a few authors. Qi et al. (2018) and Liao et al. (2019) analyzed the influence of step rotation angles on the flow patterns by numerical simulation. The results showed that with the increase in the step rotation angles, the maximum water depth on the step initially decreases and then increases, but the location of the maximum water depth is opposite. Shen et al. (2019) studied the influence of end sills on the energy dissipation of the stepped dropshaft. The results showed that the end sills on the steps can increase the energy dissipation but decrease the discharge capacity of the stepped dropshaft.

The standing wave, as an important flow phenomenon, is associated with the discharge capacity. When the height of the standing wave is too large to touch the bottom of the upper step, the discharge capacity of the stepped dropshaft will be limited. Therefore, it is necessary to control
the height of the standing wave and improve the discharge capacity. In this study, a new design of stepped dropshaft is proposed, namely the blade-shaped stepped dropshaft (Figure 1(b)). By modifying the vertical wall of the step to a curved vertical wall, it generates a more uniform flow rate distribution along the radial direction, therefore controlling the standing wave and improving the discharge capacity. ANSYS CFX 17.0 was used to simulate the flow in the stepped dropshaft. The main objectives of this study are: (1) to obtain the hydraulic characteristics of the blade-shaped stepped dropshaft in respect of flow pattern, flow rate distribution, time-averaged pressure and energy dissipation; (2) to compare with the standard stepped dropshaft and obtain the influence of the step shape on the hydraulic characteristics.

**NUMERICAL SIMULATION**

**Model setup**

Figure 1 is the definition sketch of the stepped dropshaft, which consisted of a vertical circular shaft with a tangential horizontal inlet channel and outlet channel. Air can be released through the air holes arranged on the internal wall below the steps. For both the standard and blade-shaped stepped dropshafts, the height (H) of the dropshaft, the radius (R) of the external wall, the radius (r) of the internal wall, the height (b) and step rotation angle (θ) of each step were the same: H = 1.5 m, R = 0.5 m, r = 0.2 m, b = 0.1 m, and θ = 60°. Compared with the standard stepped dropshaft, the vertical wall of the step was modified to be a curved vertical wall for the blade-shaped stepped dropshaft. The step rotation angle (α) and the radius (r1) of the curved vertical wall were α = 86° and r1 = 0.25 m, respectively. The length of the inlet and outlet channels was 1 m and 1.5 m, respectively. The cross section of the inlet and outlet channels was rectangular with 0.3 m width and 0.5 m height, respectively.

**Volume of fluid method**

The volume of fluid method (VOF) model is suitable to model the free-surface flow and track the volume fraction of each fluid over the entire region (Hirt & Nichols 1981; Qian et al. 2019). The sum of the volume fractions of water and gas in each of the calculated grid cells is constant, and they are expressed as αw and αa, respectively:

\[ α_w + α_a = 1 \]  

(1)

The density ρ and the viscosity μ can be expressed as:

\[ ρ = α_w ρ_w + (1 - α_w) ρ_a \]  

(2)

\[ μ = α_w μ_w + (1 - α_w) μ_a \]  

(3)

where ρw and ρa are the density of water and gas, respectively; μw and μa are the viscosity of water and gas, respectively.

The variables and their attributes represent air or water, or a mixture of them at any control volume. The tracking of the interface between air and water is accomplished by the continuity equation as follows:

\[ \frac{∂α_w}{∂t} + u_i \frac{∂α_w}{∂x_i} = 0 \]  

(4)

where xi and ui are the coordinates and velocity components, respectively (i = 1, 2, 3), t is the time.

**Turbulence model**

The shear stress transport rotation-curvature correction (SST-CC) model can provide an accurate prediction of the swirling flow evolution and the boundary layer simulations, which are relevant to the flow characteristics in the stepped dropshaft (Menter 1994). Many papers have simulated the complicated hydraulic properties related to the swirling flow and high accuracy boundary layer using the SST-CC model (Guo et al. 2018). Therefore, the SST-CC model can be used to model the flow in the stepped dropshaft. The governing equations are as follows:

\[ μ_T = \frac{k}{ω} \]  

(5)

\[ \frac{∂ρk}{∂t} + ρ \frac{∂(μ_t k)}{∂x_i} = P_k f_1 - β ρ \omega k + \frac{∂}{∂x_i} \left( μ + \frac{μ_T}{σ_k} \frac{∂k}{∂x_i} \right) \]  

(6)
\[
\frac{\partial \rho \omega}{\partial t} + \frac{\partial (\rho u_i \omega)}{\partial x_i} = \frac{\rho P_a}{\mu_T} f_{r1} - \rho \beta \omega^2 + 2(1 - f) \frac{\rho \sigma_{\omega \omega}}{\omega} \frac{\partial \omega}{\partial x_j} \frac{\partial x_j}{\partial x_i} + \frac{\partial}{\partial x_i} \left( \mu + \frac{\mu_T}{\sigma_\omega} \right) \frac{\partial \omega}{\partial x_i} \frac{\partial x_j}{\partial x_j} \tag{7}
\]

\[
f_{r1} = \tanh \left\{ \min \left[ \max \left( \frac{\sqrt{k}}{0.09 \nu_d^e}, \frac{500 \mu}{\rho \nu_d^e} \right), \frac{4 \sigma_\omega \omega k}{CD_{\omega \omega} d^2} \right] \right\} \tag{8}
\]

\[
CD_{\omega \omega} = \max \left( 2 \frac{\sigma_{\omega \omega}}{\omega} \frac{\partial \omega}{\partial x_j} \frac{\partial x_j}{\partial x_j}, 10^{-20} \right) \tag{9}
\]

where \( \omega \) = specific dissipation rate, \( k \) = turbulent kinetic energy, \( \mu_T \) = turbulent eddy viscosity, \( d \) = the distance to the nearest wall, and \( P_a \) and \( P_i \) = production terms. \( \beta, \beta^*, \sigma_{\omega}, \sigma_{\omega}, \sigma_{\omega \omega} \) and \( \sigma_{\omega \omega} \) are the model coefficients. The curvature correction term \( f_{r1} \) can increase the sensitivity to the curvature and rotation of the streamline, which is more in line with the spiral bending flow of the water flow inside the dropshaft (Spalart & Shur 1997). The empirical formula \( f_{rotation} \) for curvature correction is:

\[
f_{rotation} = (1 + C_{r1}) \frac{2 \pi}{1 + r^*} \left[ 1 - C_{r3} \tan^{-1} (C_{r2} \pi) - C_{r1} \right] \tag{10}
\]

\[
f_{r1} = \max \left[ \min \left( f_{rotation}, 1.25 \right), 0 \right] \tag{11}
\]

where the constant parameters are \( C_{r1} = 1.0, C_{r2} = 2.0, C_{r3} = 1.0, r^* \) and \( \tilde{r} \) are related to the strain rate tensor \( \buffer{S} \) and the magnitude of vorticity \( \Omega \):

\[
r^* = \frac{S}{\Omega} \tag{12}
\]

\[
\tilde{r} = 2 \Omega h S h k \frac{DS_h}{D_t} + \epsilon_{i} \epsilon_{mn} S_{im} + \epsilon_{j} \epsilon_{mn} S_{jm} \Omega_{mn} \left( \frac{1}{\Omega D^3} \right) \tag{13}
\]

\[
D = \max \left( S^2, 0.09 \rho^2 \right) \tag{14}
\]

**Numerical algorithm**

The Geo-Reconstruct can obtain the face fluxes whenever a cell is filled with water or air and be used when the cell is near the interface between two phases (Cui et al. 2020). It represents the interface between fluids using a piecewise-linear approach. The Pressure-Implicit with Splitting of Operators (PISO) algorithm, based on the higher degree of the approximate relation between the corrections for pressure and velocity, can decrease the number of iterations required for convergence, especially for transient problems (Xie & Xiao 2017).

**Model grid, boundary and initial conditions**

The mesh of the numerical model was configured as a structural grid with 2290037 hexahedral grid cells. Mesh refinement was used near the stepped dropshaft wall to deal with the swirling flow attached to the wall. The minimum grid size near the dropshaft wall was 2 mm. Mesh convergence was tested using a coarser mesh with 1648431 cells, showing less than 5% difference to the predicted average velocity and time-averaged pressure at the end of the eleventh step.

The upstream inlet was assumed to be the mass flow bulk. The downstream outlet was assumed to be the pressure outlet and a static pressure equal to zero was assumed over it. The top of the central hollow was assumed to be the opening boundary condition with a relative pressure equal to zero, which allows the air to cross the boundary either into or out of the domain. The wall boundary was assumed to be the no-slip wall.

The dimensionless flow rate is \( Q^* = Q/(g(R - h)^{5/2}) \), where \( Q \) is the flow rate, \( g \) is the acceleration of gravity. Numerical predictions were carried out for the range of the flow rates was \( Q = 0.012 \text{--} 0.048 \text{ m}^3/\text{s} \), which corresponded to \( Q^* = 0.078 \text{--} 0.311 \). The minimum time step was 0.001s and the max iterations per time step were set as 20.

**Model validation**

The experiment results (Wu et al. 2018) were used to evaluate the simulated settings to verify the accuracy of the calculation. In the experiment dropshaft, \( b = 0.131 \text{ m}, \theta = 150^\circ, R = 0.25 \text{ m}, r = 0.1 \text{ m}, Q^* = 0.13 \) and 0.42. Figure 2 shows the time-averaged pressure \( p \) on the ninth step horizontal surface in the flow direction and the radial direction, where \( x \) is the position of the measuring points from the vertical wall along the flow direction; \( L_x \) is the length from the vertical wall along the arranged line of the measuring points; \( y \) is the location of the measuring points.
from the internal wall along the radial direction; and \( L_Y \) is the length from the internal wall to the external wall along the radial direction.

Figure 2 showed that the maximum relative deviations in the flow and radial direction are 4.30% and 2.04%, respectively. This indicated that the results of the numerical simulation agree well with the experiment.

RESULTS AND DISCUSSION

Flow pattern

Under the influence of centrifugal force, the flow behaviors on each step of the stepped dropshaft vary in the radial direction, with the main flow remaining near the external wall. According to the variations of the main flow, referring to the flow regime classification for the stepped spillway, the flow patterns of the stepped dropshaft can be classified into nappe and skimming flows. The nappe flow is characterized by a succession of falling jets, the presence of air cavities beneath them, and hydraulic jumps downstream of the falling jets. The energy of the nappe flow is mainly dissipated by the impact of the falling jets, and the turbulence of the hydraulic jump. The characteristic of the skimming flow is that the flow skims over the steps as a coherent stream and is cushioned by the recirculating fluid trapped on the steps. The energy of the skimming flow is mainly dissipated by the intense turbulent momentum exchange between the mainstream and the recirculation.

Figure 3 shows the flow patterns of both stepped dropshafts at \( Q^* = 0.078 \) and 0.155. For the standard stepped dropshaft, the nappe flow exists at \( Q^* = 0.078 \) and 0.155. For the blade-shaped stepped dropshaft, the nappe flow exists at \( Q^* = 0.078 \) and the skimming flow exists at \( Q^* = 0.155 \). Therefore, the flow rate required for the skimming
flow in the blade-shaped stepped is significantly less than that in the standard stepped dropshaft.

Figure 3(a) and 3(b) highlight similar nappe flow performance in both dropshafts. Different from the stepped spillway, the falling jets impact on the step near the external wall under the effect of the centrifugal force, and standing waves are generated due to the deflection of the falling jets. For the standard stepped dropshaft, with the increase of $Q^*$, the jet length of the nappe flow increases gradually, resulting in the jet flow jumping over the adjacent step and landing on the lower step (Figure 3(c)). For the blade-shaped stepped dropshaft, $Q^* = 0.155$ is large enough for generating the skimming flow, with the presence of the water surface streamline parallel to the pseudo-bottom and the absence of standing waves and air cavities (Figure 3(d)).

**Flow rate distribution**

Figure 3 shows that the flow is fully developed from the eighth step at $Q^* = 0.155$. Figure 4 shows the flow rate distribution $q/Q$ with $y/L_Y$ in the cross sections at the end of the eighth to the eleventh steps at $Q^* = 0.155$, where $q$ is the flow rate integration along the water depth. For both stepped dropshafts, $q/Q$ increases from internal to external owing to the influence of centrifugal force and decreases near the external wall owing to the sidewall resistance. The maximum $q/Q$ is 0.231 and 0.112 at $y/L_Y = 0.95$ for the standard stepped dropshaft and the blade-shaped stepped dropshaft, respectively. For the standard stepped dropshaft, with increasing $y/L_Y$, $q/Q$ rises slowly at $0 < y/L_Y < 0.5$ but surges at $0.5 < y/L_Y < 0.95$. For the blade-shaped stepped dropshaft, $q/Q$ rises more slowly as a cross wavy line with increasing $y/L_Y$; therefore, the distribution of $q/Q$ is more uniform with the maximum about half of that in the standard stepped dropshaft.

It can be found that the distribution of $q/Q$ shows periodic on different steps. $q/Q$ on the eighth and tenth steps are similar, and the $q/Q$ on the ninth and eleventh steps are similar. For the standard stepped dropshaft, $q/Q$ of the eighth and tenth steps are greater than that of the ninth and eleventh steps at $0.5 < y/L_Y < 1$, which is opposite at $0 < y/L_Y < 0.5$. For the blade-shaped stepped dropshaft, $q/Q$ rises as a cross wavy line shape: when $0 < y/L_Y < 0.2$ and $0.5 < y/L_Y < 0.8$, $q/Q$ of the eighth and tenth steps are greater than that of the ninth and eleventh steps; when $0.2 < y/L_Y < 0.5$ and $0.8 < y/L_Y < 1$, it is just the opposite.

Figure 5 shows the flow rate difference $\Delta q/Q$ between the internal wall and external wall, where $\Delta q$ is the average flow rate difference between the internal wall and external wall from the eighth to eleventh steps. It shows that the relations between $\Delta q/Q$ and $Q^*$ are linear in both dropshafts. Due to the influence of centrifugal force, the $\Delta q/Q$ gradually increases with the increase of $Q^*$. Obviously, $\Delta q/Q$ of the blade-shaped stepped dropshaft is much less than that of the standard stepped dropshaft. The spacing between the two dotted lines remains approximately 0.08 independent of $Q^*$. Accordingly, the flow rate distribution of the blade-shaped stepped
dropshaft is more uniform than that of the standard stepped dropshaft. The reason will be described below via the velocity distribution on the step horizontal surface.

Velocity distribution

Figure 6 shows the velocity distribution on the horizontal step surface for the eighth to the eleventh steps in both dropshafts at $Q^* = 0.155$. It presents that after the falling jets impact the step, the water breaks up and flows in all directions: one part flows downstream directly; one part flows to the step center; the other first flows in reverse to hit the step vertical wall, then travels to the internal owing to the guidance of the step vertical wall, and finally flows downstream along the internal wall.

The quantity and location of the flow impact points (red circle in Figure 6) can be determined according to the velocity distribution. The curved vertical wall of the step first affects the number and distribution of the impact points, and then affects the distribution of $q/Q$. For the standard stepped dropshaft, influenced by the centrifugal force, the flow impact points on the step horizontal surface are mainly near the external wall, shown in Figure 6(a). For the blade-shaped stepped dropshaft, several impact points exist on the step center apart from those near the external wall, which leads the flow to the step center and the internal wall, shown in Figure 6(b). Therefore, $\Delta q/Q$ of blade-shaped stepped dropshaft is less than that of the standard stepped dropshaft (Figure 5).

Figure 5 | The flow rate difference between the internal wall and external wall.

Figure 6 | Velocity distribution on the step horizontal surface at $Q^* = 0.155$: (a) standard stepped dropshaft; (b) blade-shaped stepped dropshaft.
Time-averaged pressure

Figure 7 presents the time-averaged pressure distribution on the step horizontal surface for the eighth to eleventh steps at $Q^* = 0.078$ and 0.155. The maximum pressure on the step horizontal surface is near the external wall influenced by the impact of main flow. The pressure gradually decreases from the external wall to the internal wall in the radial direction, with a negative pressure zone near the end of the step. It indicates that the time-averaged pressure difference between internal and external increases with the increase of $Q^*$. 

![Figure 7](image_url)
Figure 7(a) and 7(b) present that due to the falling jets step by step, the time-averaged pressure distribution is similar on each step in both stepped dropshafts at $Q^* = 0.078$. Figure 7(c) presents that thanks to the same reason, with the periodic of flow rate distribution, the time-averaged pressure distribution in the blade-shaped stepped dropshaft appears periodic at $Q^* = 0.155$: pressure distribution on the eighth and tenth steps are similar, pressure distribution on the ninth and eleventh steps are similar. Figure 7(d) shows that the time-averaged pressure distribution in the blade-shaped stepped dropshaft is similar on each step, which resulted from the flow skimming over the steps as a coherent stream.

There is an obvious negative pressure zone on each step vertical wall. Figure 8 presents the dimensionless pressure distribution $p_n/\rho g b$ on the negative pressure zone with $y/L_Y$ for the eighth to eleventh steps at $Q^* = 0.155$, where $p_n$ is the vertical average pressure along the water depth. With increasing $y/L_Y$, $p_n/\rho g b$ first decreases under the influence of the centrifugal force, and then increases near the external wall due to the sidewall resistance. Corresponding to the time-averaged pressure distribution on the step horizontal surface, $p_n/\rho g b$ shows periodic for the standard stepped dropshaft: $p_n/\rho g b$ of the eighth and tenth steps are similar with the minimum $p_n/\rho g b$ at $y/L_Y = 0.67$; $p_n/\rho g b$ of the ninth and eleventh steps are similar with the minimum $p_n/\rho g b$ at $y/L_Y = 0.93$ (Figure 8(a)). For the blade-shaped stepped dropshaft, $p_n/\rho g b$ from the ninth to eleventh steps are similar, with the minimum $p_n/\rho g b$ at $y/L_Y = 0.68$.

Figure 9 shows the minimum $p_n/\rho g b$ and the dimensionless negative pressure area $s/lb$, where $s$ is the average negative pressure area for the eighth to eleventh steps vertical walls. It indicates that both $p_n/\rho g b$ and $s/lb$ linearly decrease with the increase of $Q^*$. With the blade-shaped stepped dropshaft, the minimum $p_n/\rho g b$ are greater and the $s/lb$ are less than that of the standard stepped dropshaft.

**Energy dissipation**

The energy dissipation $\eta$ of the dropshaft is expressed as follows:

$$\eta = \frac{E_1 - E_2}{E_1}$$  \hspace{1cm} (15)

$$E_1 = Z_1 + \frac{P_1}{\rho g} + \frac{v_1^2}{2g}$$ \hspace{1cm} (16)

$$E_2 = Z_2 + \frac{P_2}{\rho g} + \frac{v_2^2}{2g}$$ \hspace{1cm} (17)

where $E_1$ and $E_2$ are the total energy at the inflow and outflow sections, respectively; $Z_1$ and $Z_2$ are the elevation heads of the inlet and outlet channels, respectively; $P_1/\rho g = 0$ and $P_2/\rho g = 0$ are the pressure heads of the inlet and outlet channels, respectively; $v_1$ and $v_2$ are the velocity heads at the inflow and outflow sections, respectively.
and outlet channels, respectively; $v_1$ and $v_2$ are the average velocities of the inflow and outflow sections, respectively.

Figure 10 shows the energy dissipation in various dropshafts. The energy dissipation decreases with the increase in the $Q^*$. The energy dissipation of the standard stepped dropshaft is similar to that of the blade-shaped stepped dropshaft, which indicates that the curved vertical wall of the step has little effect on the energy dissipation. $\eta$ of the stepped dropshaft in this study agree well with experimental data of the standard stepped dropshaft from Wang (2018), and the expression of the fitting curve is obtained as:

$$
\eta = 0.6689(Q^*)^{-0.127}
$$

A comparison with the results of a plunging dropshaft (Chanson 2004) indicates that $\eta$ of the stepped dropshaft is similar to that of the plunging dropshaft in some cases. It is believed that $\eta$ is related to the multiple parameters of the dropshaft, which is needed further studies.

**CONCLUSIONS**

In order to control the standing wave and improve the discharge capacity, the blade-shaped stepped dropshaft was proposed in this study to generate a more uniform flow rate distribution in the radial direction by modifying the vertical wall to a curved vertical wall. Three-dimensional numerical simulation was carried out for the hydraulic characteristics of the blade-shaped stepped dropshaft in respect of flow pattern, flow rate distribution, time-averaged pressure and energy dissipation, and the results are compared with the standard stepped dropshaft. The conclusions are summarized as follows:

Compared with the standard stepped dropshaft, the flow rate distribution of the blade-shaped stepped dropshaft was more uniform, with less flow rate difference between the
internal and external walls at each discharge. Due to the curved vertical wall of the step, more impact points appear on the center of the horizontal step surface, which lead the flow to the internal wall against the centrifugal force. By improving the flow rate distribution in the radial direction, the flow concentration on the external wall and the standing wave can be well controlled.

The numerical simulation results showed that there is a periodic phenomenon of the flow rate distribution and pressure distribution in the stepped dropshaft, which resulted from the flow jumping over the adjacent step and landing on the lower step. Compared with the standard stepped dropshaft, the blade-shaped stepped dropshaft has more uniform pressure distribution on the horizontal step surface and more limited negative pressure on the vertical wall owing to the adjustment of the numbers and locations of the impact points. The energy dissipation of the blade-shaped stepped dropshaft is as high as that of the standard stepped dropshaft, which remains above 0.7 within the range of flow rates.

Based on these findings, it can be found that compared with the standard stepped dropshaft, the blade-shaped stepped dropshaft can improve the discharge capacity, improve the pressure distribution, and keep a good energy dissipation. Therefore, it is a potentially safe and efficient dropshaft design that can be applied to deep tunnel systems.

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**DATA AVAILABILITY STATEMENT**

All relevant data are included in the paper or its Supplementary Information.

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