COOL-LAMPS. I. An Extraordinarily Bright Lensed Galaxy at Redshift 5.04*

Gourav Khullar1,2, Katya Gozman1,3, Jason J. Lin1,2, Michael N. Martinez1,3, Owen S. Matthews Acuña1,2, Elisabeth Medina1, Kaiya Merz1, Jorge A. Sanchez1, Emily E. Sisco1,2, Daniel J. Kavin Stein1, Ezra O. Sukay1, Kiyano Tavangar1, Matthew B. Bayliss4, Lindsey E. Bleem2,5,6,7, Sasha Brownberger6,8, Hakon Dahle7,8, Michael K. Florian9, Michael D. Gladders1,2,9, Guillaume Mahler3, Jane R. Rigby3, Keren Sharon3, and Antony A. Stark9

1 Department of Astronomy and Astrophysics, University of Chicago, 5640 South Ellis Avenue, Chicago, IL 60637, USA; gkhullar@uchicago.edu
2 Kavli Institute for Cosmological Physics, University of Chicago, 5640 South Ellis Avenue, Chicago, IL 60637, USA
3 Department of Astronomy, University of Michigan, 1085 South University Drive, Ann Arbor, MI 48109, USA
4 Department of Physics, University of Cincinnati, Cincinnati, OH 45221, USA
5 Argonne National Laboratory, High-Energy Physics Division, 9700 S. Cass Avenue, Argonne, IL 60439, USA
6 Department of Physics, Harvard University, 17 Oxford Street, Cambridge, MA 02138, USA
7 Institute of Theoretical Astrophysics, University of Oslo, P.O. Box 1029, Blindern, NO-0315 Oslo, Norway
8 Observational Cosmology Lab Code 665, NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA
9 Center for Astrophysics | Harvard & Smithsonian, 60 Garden St, Cambridge, MA 02138, USA

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Abstract

We report the discovery of COOL J1241+2219, a strongly lensed galaxy at redshift $z = 5.043 \pm 0.002$ with observed magnitude $z_{AB} = 20.47$, lensed by a moderate-mass galaxy cluster at $z = 1.001 \pm 0.001$. COOL J1241 +2219 is the brightest lensed galaxy currently known at optical and near-infrared wavelengths at $z \gtrsim 5$; it is $\sim 5$ times brighter than the prior record-holder lensed galaxy, and several magnitudes brighter than the brightest unlensed galaxies known at these redshifts. It was discovered as part of COOL-LAMPS, a collaboration initiated to find strongly lensed systems in recent public optical imaging data. We characterize the lensed galaxy, as well as the central galaxy of the lensing cluster using ground-based $grizJH$ imaging and optical spectroscopy. We report model-based magnitudes, and derive stellar masses, dust content, metallicity, and star-formation rates via stellar-population synthesis modeling. Our lens mass modeling, based on ground-based imaging, implies a median source magnification of $\sim 30$, which puts the stellar mass and star-formation rate (in the youngest age bin, closest to the epoch of observation) at $\log M_\star = 10.11^{+0.23}_{-0.28}$ and $\text{SFR} = 27^{+13}_{-9} M_\odot \text{yr}^{-1}$, respectively. We constrain a star-formation history for COOL J1241+2219 consistent with constant star formation across $\sim 1$ Gyr of cosmic time, and that places this galaxy on the high-mass end of the star-forming main sequence. COOL J1241+2219 is two to four times more luminous than a galaxy with the characteristic UV luminosity at these redshifts. The UV continuum slope $\beta = -2.2 \pm 0.2$ places this galaxy on the blue side of the observed distribution of galaxies at $z = 5$, although the lack of Ly$\alpha$ emission indicates dust sufficient to suppress this emission.

Unified Astronomy Thesaurus concepts: High-redshift galaxies (734); Galaxies (573); Ultraviolet astronomy (1736); Astronomical methods (1043); Star formation (1569); Gravitational lensing (670); Strong gravitational lensing (1643); Surveys (1671); Spectral energy distribution (2129); Photometry (1234); Spectroscopy (1558); Galaxy properties (615)

1. Introduction

Strong gravitational lensing can transform typically small, faint, marginally resolved field galaxies into highly magnified arcs that can appear very bright. The brightest and most magnified lensed galaxies enable exquisite studies of internal galaxy morphologies down to scales of only tens of parsecs even in the most distant galaxies (e.g., Bayliss et al. 2014; Livermore et al. 2015; Johnson et al. 2017a, 2017b; Cornachione et al. 2018; Rivera-Thorsen et al. 2019; Florian et al. 2020; Ivison et al. 2020). Such lensed sources are exceedingly rare, and are generally found by searching many thousands of square degrees of ground-based survey imaging (e.g., Alam et al. 2007; Belokurov et al. 2007; Koester et al. 2010) or by dedicated imaging programs targeting massive galaxy clusters not previously explored (Bleem et al. 2015, 2020; Cee et al. 2019). Conversely, deep observations of known strong lenses with large lensing cross-sections (see the Hubble Frontier Fields; Lotz et al. 2017) discover many more lensed galaxies, but these are magnitudes fainter than the rare systems found in wide searches.

Studies of the brightest lensed galaxies are complementary to studies of faint lensed galaxies, and unlensed field galaxies, offering both enhanced spatial resolution in the source plane, and sufficient photons to allow exceptional spectral resolution at useful signal-to-noise ratios. Hence, the brightest strongly lensed galaxies are the observational signposts with which we navigate the more abundant but much fainter deep-field samples. To date, searches of the Sloan Digital Sky Survey (SDSS; Ahumada et al. 2020) have identified hundreds of bright lensed sources at $z < 3$, the brightest of which typically have integrated $g_{AB}$ magnitudes of $\sim 20$ (e.g., Kubo et al. 2010; Bayliss et al. 2011; Stark et al. 2013). This success is due to the depth and filter selection of the SDSS, sufficient to find the bright blue arcs that are the hallmark of UV-bright lensed galaxies at these redshifts. Many of these lensed targets have been followed up with HST imaging that has revealed details of star formation and structure

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* This paper includes data gathered with the 6.5 m Magellan Telescopes located at Las Campanas Observatory, Chile.
on spatial scales down to tens of parsecs (e.g., Johnson et al. 2017a; Cornachione et al. 2018). The brightest of these lensed galaxies, a few of which predate the discoveries from the SDSS, have been followed up with high-quality spectroscopy (e.g., Pettini et al. 2000; James et al. 2014; Rigby et al. 2018a, 2018b) that exceeds in both spectral resolution and signal-to-noise what is available for stacks of dozens or even hundreds of field galaxies (Shapley et al. 2003; Steidel et al. 2016).

Beyond $z \sim 3$ however, bright lensed samples rapidly decline, and by $z \sim 5$, there are only a few known strongly lensed galaxies of even modest brightness at optical/infrared (OIR) wavelengths; the SDSS, which has been so effective for discovery at $z \sim 2$–$3$ has proven mostly insufficient in depths at $z \sim 5$, particularly in key redder filters as the galaxies drop out of the bluer filters due to intervening IGM absorption. The brightest three $z \sim 5$ spectroscopically confirmed lensed galaxies published to date have AB magnitudes of $i = 22.2$ ($z = 4.88$, Gladders et al. 2002), $i = 22.6$ ($z = 4.92$, Franx et al. 1997; Soifer et al. 1998), and $i = 23.3$ ($z = 4.87$, Frye et al. 2002). The brightest was found through a wide-area search for arcs; the other two were found by targeting massive galaxy clusters. In addition, the Reionizing Lensing Cluster Survey (RELICS) program reports five candidate lensed galaxies with $z_{\text{phot}} = 5.5$–$6$ and $F160W < 23.6$ (Salmon et al. 2020), that as yet lack published spectroscopic redshifts. For comparison, the brightest galaxies at $z \sim 5$ in the Hubble Ultra Deep Field (HUDF) have magnitudes fainter than $i_{\text{AB}} = 25$ (Inami et al. 2017).

Strong-lensing searches in the latest generation of imaging surveys, particularly the Dark Energy Camera Legacy Survey (DECaLS; Huang et al. 2020), the Dark Energy Survey (DES; Diehl et al. 2017; Jacobs et al. 2019), the Kilo-Degree Survey (KIDS; Petrillo et al. 2019), the Hyper Suprime-Cam Subaru Strategic Program (HSC-SSP) Survey (Jelaini et al. 2020), should significantly increase the number of known bright $z > 3$ lensed galaxies, as these data are deeper, particularly in redder filters, than the SDSS. Here, we report and characterize one such object, a spectroscopically confirmed $z = 5.043$ strongly lensed galaxy, with an extraordinarily bright $i_{\text{AB}}$ apparent magnitude of 20.47. As such, this is the brightest $z \gtrsim 5$ UV-bright galaxy found so far.

This paper is structured as follows. Section 2 lays out the details of the public data search that led to the discovery of COOL J1241+2219. Section 3 describes follow-up imaging and spectroscopy of COOL J1241+2219, while Section 4 describes the analyses of these data and the results from stellar-population synthesis and strong lens modeling. Section 5 places these results in the context of other objects known in the $z > 5$ universe.

All reported magnitudes are calibrated to the AB system. The fiducial cosmology model used assumes a standard flat cold dark-matter universe with a cosmological constant ($\Lambda$ CDM), corresponding to WMAP9 observations (Hinshaw et al. 2013). For inferred parameters with uncertainties, we report 16th, 50th, and 84th percentile values, unless otherwise specified.

2. Discovery

The COOL-LAMPS projectChicagO Optically selected strong Lenses—Located At the Margins of Public Surveys initiated as the central focus of an undergraduate research class, is an effort to find strong gravitational lenses in recent public imaging data. Though the lensing search is designed to find a wide variety of lenses and lensed sources, the particular focus of our follow-up is sources that are photometrically at the margins of the distributions of source color and brightness. The details and results of this search will be presented in an upcoming publication (COOL-LAMPS Collaboration 2020, in preparation), and we provide only a brief description here.

COOL J1241+2219 was found in the DECaLS grz Data Release 8 (DR8, Dey et al. 2019) imaging data, as part of a comprehensive visual search of the northern galactic cap portion of the DECaLS data set. The complete search comprised two separate efforts, one targeting potential lenses at low redshift ($z < 0.7$) and the other extending to higher redshift potential lenses. COOL J1241+2219 was found in the latter.

We selected luminous red galaxies, classified as having de Vaucouleurs profiles (de Vaucouleurs 1948), and appropriate colors and magnitudes, as potential lenses. Color and magnitude cuts were informed by prior experience with strong lens selection in the SDSS (Bayliss et al. 2011; Sharon et al. 2020); to first order, the selected galaxies are objects brighter and redder than a color–magnitude model corresponding to a passively evolving galaxy formed at $z = 3$, with an absolute magnitude tuned to capture a reasonable number of target objects. In practice, this magnitude limit is approximately $M_{ \text{g}}$. Specifically, COOL J1241+2219 ’s measured photometry puts it at $\sim 0.2$ mag redder and brighter than the selection cut, and further away from the selection cut than the median potential lens selected by that cut.

We followed a visual examination strategy employed by the Sloan Giant Arcs Survey (SGAS; e.g., Bayliss et al. 2011; Sharon et al. 2020). In the high-redshift portion of our search, a total of 271,460 lines of sight were examined by at least three of the coauthors, primarily via a custom Python-based image viewer and lens-ranker. We examined lines of sight four at a time in a $2 \times 2$ grid, with each individual cutout covering $45'' \times 45''$ centered on each potential lens. We constructed custom color images from the DECaLS grz imaging data, tuned to emphasize faint extended features. Each $2 \times 2$ composite image was given an integer score from 0 to 3 by each examiner, with a score of 3 indicating a definite strong lens, a score of 0 indicating no evidence for strong lensing, and the remainder shading between those extremes. The initial average score for COOL J1241+2219 was 2. The central portion of the image in which COOL J1241+2219 was found is shown in the first inset panel of Figure 1.

The lensed arc is detected in the DECaLS catalogs at R.A., decl. (J2000) = (190.3743, 22.3282) with reported AB magnitudes of $g = 25.10$, $r = 23.82$, $z = 21.13$; the model-based photometry of these catalogs assumes elliptical sources, and must therefore miss a significant fraction of the source light in this case. However, this initial photometry and the observed red color and morphology of this apparent arc motivated follow-up observations.

3. Follow-up Observations and Redshifts

3.1. Photometry

Near-infrared imaging of COOL J1241+2219 in the J- and H-bands was obtained using the FourStar Infrared Camera...
on the Magellan/Baade telescope, Chile, on 2020 February 22. The total integration times were 335 s and 3578 s in the $J$- and $H$-bands, respectively. Data were reduced to final astrometrically and photometrically calibrated stacked images using a custom pipeline built via IRAF and PHOTPIPE (Rest et al. 2005; Garg et al. 2007; Miknaitis et al. 2007).

We then obtained deep simultaneous optical imaging—griz—using the Parallel Imager for Southern Cosmology Observations (PISCO; Stalder et al. 2014) on the Magellan/Clay telescope, on 2020 February 29. The total integration time for the imaging with a field of view of $5.5 \times 7.5$ was 3060 s. Data were processed and stacked to final astrometrically calibrated images using a custom pipeline that makes standard corrections, as well as removing effects of nonlinearity and bright stars peculiar to the PISCO detectors. Figure 1 shows a color image for COOL J1241+2219 constructed from the PISCO images, as well as smaller grayscale cutouts for all of the follow-up imaging focused on the arc and central lens galaxy.

As can be seen in Figure 1, these data show that COOL J1241+2219 is blue in the $z$–J and $z$–H colors. In combination with the observed red optical colors this indicates that this arc is the strongly lensed image of a high-redshift galaxy.

3.2. Spectroscopy

Spectroscopic observations for COOL J1241+2219 were obtained on 2020 March 2, using the Low Dispersion Survey Spectrograph—3C (LDSS-3C\(^{10}\)), with the VPH-Red grism and a 1\%'0 long slit placed to cover a wavelength range of 6700 Å and redward, to the limits of the instrument response at $\sim$10500 Å. The slit was oriented to include the brightest clump in the lensed arc, the fiducial counterimage, and the apparent lens galaxy (with a slight offset). The precise slit placement is shown in the top inset portion of Figure 2.

We acquired six 900 s integrations, dithered between two slit positions along the slit direction, in non-photometric conditions at airmasses of 1.6–1.8, with a typical on-instrument seeing of 1\%'1, as well as observations of a spectrophotometric standard for both flux calibration and the removal of telluric features in the spectra. Reduction to wavelength- and flux-calibrated 1D spectra was accomplished using a combination of standard routines and custom software, implemented in IRAF and IDL. The spectra for the apparent central lens galaxy and COOL J1241+2219 are shown in Figure 2, along with an estimate of the noise spectrum, and the applied telluric correction. We

\(^{10}\) http://www.lco.cl/?epkb_post_type_1=ldss-3-user-manual
measure the lens redshift from the marked features in Figure 2 as $z = 1.001 \pm 0.001$ and the lensed source redshift from cross-correlation with the stacked rest-UV galaxy spectrum from Rigby et al. (2018a) as $z = 5.043 \pm 0.002$.

The high-airmass and non-parallactic slit positioning of these observations in suboptimal conditions makes the fluxing and telluric calibrations irretrievably uncertain. However, as is obvious from Figure 2, both spectra are sufficient to establish redshifts robustly.

At the measured redshift, the entire rest-frame optical band for the lensed galaxies is shifted to wavelengths not readily accessed from ground-based observatories. The only observable standard emission line used in studies of galaxies at lower redshifts is the $[\text{O II}]_{3727,3729}$ Å doublet. We observed COOL J1241+2219 using the Gemini-North Near-infrared Spectrograph (GNIRS) spectrograph (Elias et al. 2006) at NOIRLAB/Gemini, Maunakea, in Director’s Discretionary Time on 2020 July 4 in an attempt to measure this feature. We acquired eight 257 s integrations using GNIRS with a 0"675 slit in cross-dispersed mode and with the short camera, at an airmass~1.5, with a typical on-instrument seeing of 0"6. These observations targeted the same bright peak along the arc observed with LDSS-3C. The data were reduced primarily using the Gemini/GNIRS pipeline. The resulting spectrum was calibrated against a telluric standard star observed immediately after. Analysis of these data is detailed in Section 4.

4. Analysis and Results

4.1. Zero-point Calibration

The $J$- and $H$-band images are calibrated to 2MASS stars (Jarrett et al. 2000) within the field of view, with the calibration derived automatically by PHOTPIPE routines. Uncertainties on these zero-points relative to 2MASS are $\sim 0.03$ mag.

The PISCO data are calibrated to existing DECaLS (griz-bands) and Pan-STARRS (DR2, Chambers et al. 2016) (i-band) data. We used SExtractor (Bertin & Arnouts 1996) MAG_AUTO magnitudes for detected objects in all bands to establish image

11 https://www.gemini.edu/instrumentation/gnirs/data-reduction

Figure 2. The extracted optical/near-IR Magellan/LDSS3 spectra of the central lens galaxy (top), brightest knot in the arc (middle top), the noise spectrum (middle bottom) and telluric correction applied (bottom). The noise and transmission spectra are shown covering the same observed wavelength range, corresponding to the indicated rest-frame wavelength range for redshifts of $1.001 \pm 0.001$ and $5.043 \pm 0.002$ for the top and middle top panels, respectively. The raw spectra are plotted in pale blue, and versions of those data filtered (via a median and top-hat filter kernel) to visually emphasize features relevant for redshift determinations, are overplotted in orange. The inset at top left shows the slit positioning for these observations. In the top panel, vertical dotted lines indicate significant spectral features used to infer the redshifts for the lens galaxy; black lines (from left to right) mark the Ca II H&K lines, and the G-band. For the lensed source spectrum, we mark, for reference, Ly$\alpha$ in absorption at 1215 Å, though note that the redshift is determined from an uncertainty-weighted cross-correlation with the reference spectrum (from Rigby et al. 2018a) overplotted in green. To ease visual comparison, we have scaled the reference spectrum using a low-order polynomial computed redward of Ly$\alpha$ to match its amplitude to the lensed source spectrum. To visually confirm the redshift, we mark UV stellar and nebular absorption features with red lines (from left to right, in rest-frame Å units) $S \Pi 1260, O\ I/S \Pi 1302/1304, C \ II 1334, S \ IV 1393, S \ II 1526, C \ IV 1548/1550$. 

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zero-points. The PISCO data are much deeper than the reference data, and so object-by-object uncertainties in the comparison are dominated by uncertainties in the reference measurements. Limiting the comparison to objects with a measured DECaLS SNR of $>20$ yields a matched catalog sufficient to establish zero-points, and linear color terms, accurate to better than 0.02 mag, and we derive a similar result in comparison to the Pan-STARRS data.

### 4.2. Model Photometry with GALFIT

We used the parametric fitting code GALFIT (Peng et al. 2002, 2010) to create a model of the arc, its candidate counterimage, the central lens galaxy, and other nearby contaminating objects, as well as to extract relevant photometry. Point-spread functions (PSFs) were measured directly from the images; these are simply small cutouts of two isolated and bright but unsaturated point sources near the modeled region. Differences in photometry from choices of PSF are substantially smaller than all uncertainties below. Multiple 2D Sérsic components were used to model each image. Model construction proceeded iteratively, with bright objects fit first, and further components added until the residual was consistent with the background noise.

To estimate statistical photometric uncertainties, we added the final fitted model to blank sky regions around the image, and refit the models on these new images. The measured magnitude distribution of each object then gives an uncertainty. If the flux distribution had significant outliers (presumably due to poor convergence of the model instance), we computed the uncertainty from the flux distribution iteratively clipped at $3\sigma$. Additionally, to estimate the systematic uncertainty caused by modeling choices, three coauthors independently created models for the $z$-band data, and the range in final fitted magnitudes for these models was taken as an estimate of the systematic uncertainty for the PISCO images. Additionally these models were re-optimized to fit the $H$-band image, and we provide an estimate of the systematic uncertainty in the $J$- and $H$-bands.

Finally, calibration zero-point uncertainties and the statistical and systematic uncertainties estimated as above were combined in quadrature to compute total uncertainties for each measurement of each physical object.

The $i$-filter image, GALFIT model, and residuals for the lens field are shown in Figure 3. Table 1 gives the model magnitudes for the arc and the candidate counterimage.

### 4.2.1. Lens Galaxies

As part of the modeling, we have measured photometry for the central lens galaxy, as well as two other nearby galaxies to the east that have similar apparent colors and morphologies (and are visible in Figure 1). These two other galaxies were fit together, but independently of other objects. These two galaxies have colors similar to the apparent central lens galaxy, as is indicated also by the DECaLS photometry of these galaxies. Independent of strong-lensing mass modeling, this suggests that the lens is a lower-mass galaxy cluster.

### 4.2.2. Arc and Possible Counterimage

The lensed source is absent in the $g$-band and nearly so in the $r$-band, both due to absorption from the intergalactic medium (IGM). It is not possible to fit a model directly in the $g$-band.

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Figure 3. $12'' \times 12''$ cutouts, all at the same intensity scaling, of the (left) PISCO $i$-band image, (center) GALFIT model of the central lens galaxy, the arc, and other galaxies, and (right) the residual image (with inverted colors to highlight the lack of residual structures).

| Model Photometry-based Magnitudes for COOL J1241+2219 |
|------------------------------------------------------|
| **$g$** | **$r$** | **$i$** | **$z$** | **$J$** | **$H$** |
|----------|---------|---------|---------|-------|-------|
| Central Lens Galaxy | 23.86$^{+0.24}_{-0.20}$ | 21.960 ± 0.046 | 20.76 ± 0.05 | 19.630 ± 0.048 | 18.955 ± 0.059 | 18.564 ± 0.054 |
| Arc | $>25.11$ | 23.197$^{+0.353}_{-0.358}$ | 21.616$^{+0.049}_{-0.150}$ | 20.472$^{+0.017}_{-0.048}$ | 20.53$^{+0.273}_{-0.265}$ | 20.566$^{+0.084}_{-0.087}$ |
| Counterimage | 24.164$^{-0.233}_{-0.297}$ | 22.705$^{-0.105}_{-0.127}$ | 23.074$^{-0.453}_{-0.795}$ | 23.774$^{-1.545}_{-1.413}$ |

**Note.** Magnitudes calculated via GALFIT, with output in AB system. $griz$-band imaging from Magellan/PISCO simultaneous multi-band imager. $JH$ band imaging from Magellan/FOURSTAR infrared imager.
Rather, we took the initial $i$-band model, solved for any residual astrometric shift between the two bands using visible objects in the modeled region, and then performed a constrained refit of the model allowing only the magnitudes to vary of any components that describe the arc. This model was pushed through the same statistical error analysis as other filters, though in the case of the $g$-band image the distribution of notional measured magnitudes represents an upper limit, and we report the 95th percentile value.

In the $r$-, $J$-, and $H$-band data, the arc is visible but measured at a markedly lower SNR than the $i$- and $z$-bands. In each of these cases, we explored both direct modeling and re-optimized models taken from bands with larger SNR detections of the arc; the measurements we report come from the latter. Note that, in all of these cases, including the $g$-band, components that describe well-detected objects were allowed to vary structurally, as well as in brightness.

The suggested counterimage (see Figure 1) is photometrically consistent with the arc; it has consistent colors across filters where both objects are visible, and the apparent surface brightness of this object matches much of the arc. The candidate counterimage was observed via LDSS-3C optical long-slit spectroscopy (and not with GNIRS) but, given that the lensed source has no emission lines in the rest-frame UV, we have not attempted to extract a spectrum of this much fainter source. We note, however, that there is no evidence in the 2D spectral image of any emission lines from this object that would be inconsistent with the identification of it as a lensed counterimage.

### 4.3. SED Fitting with Prospector

Using the model photometry in $grizJH$ filters, we perform spectral energy distribution (SED) fitting for the central lens galaxy and the COOL J1241+2219 arc using the Markov Chain Monte Carlo (MCMC)-based stellar-population synthesis (SPS) and parameter inference code, Prospector. Prospector is based on the Python-FSPS framework, with the MILES stellar spectral library and the MIST set of isochrones (Conroy & Gunn 2010; Falcón-Barroso et al. 2011; Foreman-Mackey et al. 2013; Choi et al. 2016; Johnson & Leja 2017; Leja et al. 2017). We exclude the LDSS3 spectrum from the fitting procedure because, at its low short-to-noise ratio, it mostly samples saturated interstellar absorption lines that are not directly constrained by Prospector libraries.

We have chosen to perform the SED fitting on the observed photometry, a choice which assumes that there are no correlations within the SED fitting that couple parameters that are magnification-dependent (such as stellar mass, star-formation rate) with those that are not magnification-dependent (such as metallicity, dust extinction). This assumption has been tested by arbitrarily rescaling the input photometry by the measured mean magnification (see Section 4.5), and the resulting SED fits are entirely consistent with the results presented below.

In these models, we assume a nonparametric star-formation history (with age bins with [0–50], [50–100], and [100–1000] Myr in look-back time), represented by the parameters $SFR_{ratio}$ referring to the ratio of total star formation in adjacent time bins. These bins correspond to redshift bins [5.04–5.22], [5.22–5.41], and [5.41–20.50]. The priors for the star-formation-rate ratios in adjacent time bins is the continuity prior, which fits directly for the change in log(SFR) between the bins and weights against sharp transitions in SFR (see Leja et al. 2019).

We also fit for dust attenuation, using the Calzetti et al. (2000) attenuation, applied to all light from the galaxy (in units of opacity at 5500 Å), stellar metallicity $\log(Z/Z_\odot)$ (where $Z_\odot = 0.0142$), gas ionization parameter $U$, gas-phase metallicity (linked to stellar metallicity) and total mass formed in the galaxy ($M_{tot}$, in units of $M_\odot$), as free parameters. For each chain in the MCMC analysis, the remnant stellar mass $M_h,remnant$ after compensating for mass loss from post main-sequence stars, and stellar mass locked in stellar remnants. The dust extinction and metallicities have flat and liberal priors, roughly covering the range allowed by the spectral model libraries. These models each assumed a Kroupa IMF (Kroupa 2001). Nebular continuum and line emission are present, and a nominal velocity smoothing of 250 km s$^{-1}$ was used for the spectrum, informed by and consistent with both an SED fit with this as a free parameter as well as a direct fit of the spectrum using a simple Gaussian model (see below).

The arc is not visible in $g$, and the $r$- and $i$-bands sample the Ly$\alpha$ forest and Ly$\alpha$ and Ly$\beta$ in the arc. To avoid biases arising from a lack of correlation between star-formation activity and the complexity of Ly$\alpha$ absorption/emission in galaxies, which is not captured in the stellar-population synthesis models, we use only $zJH$ photometry for the arc SED fitting. We add to this the IR spectroscopy sampling the [O II] 3727, 3729 Å doublet. The spectrum and photometry are fit here simultaneously.

We assume that the [O II] equivalent width is uniform across COOL J1241+2219 in the source plane, and we scale the emission line spectrum to be representative of the whole arc by considering the entire arc flux against the portion contained within the slit, in the $i$-band. This flux normalization is uncertain in practice, quite apart from the underlying assumption of uniform equivalent width, so we include a nuisance parameter, specnorm (spectrum normalization factor), to capture any unresolved differences between model photometry and spectroscopy. This parameter is simultaneously fit with other parameters within Prospector. Note however, that this parameter is computed without this parameter in place are consistent with conclusions from models (e.g., stellar mass) that include specnorm.

The best-fit SED model is shown in Figure 4.d. We also show corner plots constructed via pyGTC (Bocquet & Carter 2016), with 1D and 2D projections of the posterior probability functions (to demonstrate both 1D marginalized probabilities and 2D covariance) in Figure 5 for the free parameters fit in the arc SED modeling (see Table 2).

It is also crucial to note that the best-fit SED model fits a nominal Gaussian model to the observed wavelength range where we expect [O II] emission. The result here (mean, width, and amplitude) of the [O II] emission feature, is consistent with results from fitting a 1D Gaussian profile to just the spectroscopic data (see Section 4.4 for details).

The parameter $M_{tot}$, total stellar mass formed, converges to a trailed Gaussian posterior distribution for these sources, which corresponds to a model-generated remnant stellar mass distribution, displayed in the bottom panel of Figure 5. The median value for the remnant stellar mass ($M_h$) in the image plane (i.e., the magnified value) is $\log(M(M_h)) = 11.63^{+0.19}_{-0.24}$. The precision of this result may surprise the reader, given that it is based on rest-UV photometry, which is, under most circumstances (e.g., at lower redshifts, and in the presence of
Figure 4. (Left) Best-fit SED model with nebular line emission (blue) and corresponding residual ($\chi$) values (green) for COOL J1241+2219, calculated via Prospector, and using zJH photometry and IR spectroscopy. Best-fit SED model without nebular line emission is displayed as short-dashed gray curve. Best-fit photometry for zJH data is shown as black squares. Photometric projections for the masked $ri$ data corresponding to the model with and without nebular line emission are shown as blue “+” and black “plus” symbols, respectively. (Right) Best-fit model spectrum with nebular line emission (blue) zoomed into the observed wavelength range corresponding to the blended [O II] 3727, 3729 Å emission, with IR spectroscopic data (orange) and uncertainties (dotted green), with model $\chi$ values (in green) displayed in the bottom panel. Two black vertical lines mark the locations of the [O II] 3727, 3729 Å doublet peaks.

significant extinction) insufficient to constrain stellar mass reliably. However, only modest dust columns are allowed by the data and, at $z = 5.043$, the universe was not old enough to conceal much stellar mass in older stars that do not produce significant UV flux.

We further test different models to check the robustness of our results. The stellar mass of the arc is found to be robust when varying parameters as follows:

1. Stellar metallicity fixed at the median best-fit value in Figure 5.
2. Removing nebular line emission from the model.
3. Altering the number of age bins from three to four, which is recommended by Leja et al. (2019) (as well as from three to two), spanning the same time period.
4. Changing mean and standard deviation of the star-formation history priors.
5. Assuming a parametric star-formation history (decayed-tau with a single burst).
6. Constraining SFH only to 0–500 Myr, as opposed to 0–1000 Myr.
7. Using velocity smoothing as a free parameter for the [O II] emission line fitting.
8. Fixing gas ionization parameter to $-2.5$, as well as gas-phase metallicity to that of the stellar value (i.e., 0).

We use the Bayesian Information Criterion (BIC) to quantify our model selection, for the greater penalty imposed by BIC on the number of free parameters used in our fitting procedure. Given a candidate model and $M$ models to compare with, the magnitude of $\Delta$BIC ($\text{BIC}_M - \text{BIC}_{\text{candidate}}$) can be inferred as evidence against a candidate model being the best model (Kass & Raftery 1995). The candidate best-fit model has a BIC value of $-1245.5$, and the models discussed above have a $|\Delta$BIC| < 5, which implies that the evidence against our candidate best-fit model is positive, albeit low. Qualitatively, all models constrain stellar mass and dust attenuation, while leaving metallicities and the gas ionization parameter unconstrained (albeit used here as nuisance parameters).

It is important to note that the SED-fitting model contains contribution by nebular line emission in order to fit [O II] emission, and hence depicts Ly$\alpha$ in emission even though it is absent in the spectrum and photometry (as indicated by mismatch between observed photometry and the synthetic magnitudes marked by plus symbol in Figure 4). This does not impact our results since the photometry sampling Ly$\alpha$ is not included in our candidate best-fit model, and $M_*$ derived from removing nebular emission contribution is within uncertainties consistent with that derived from our best-fit model.

To cross-check the complete end-to-end pipeline from raw data to SED fits, we have also applied a Prospector-based analysis to measurements of the central lens galaxy and two likely cluster members discussed in Section 4.2.1, and find that these galaxies are well fit by passively aging old stellar populations at the notional lens redshift, as expected.

4.4. Instantaneous SFR: [O II]

Distinct from the simultaneous spectrophotometric SED fit to the data, we cross-check our SFR result by calculating the [O II] 3727, 3729 Å doublet line flux to characterize the instantaneous SFR (Rosa-González et al. 2002) for COOL J1241+2219 from a scaling relation. We fit a four-parameter 1D Gaussian profile (with mean, width, amplitude, and continuum pedestal) to the spectroscopic data, in the rest-frame
region 3720–3735 Å, to capture the blended-doublet emission line. The observed [O II] line flux, corrected for the entire arc, is $4.5^{+0.4}_{-0.3} \times 10^{-16}$ erg cm$^{-2}$ s$^{-1}$. The redshift derived from blended-doublet line is $5.041 \pm 0.001$.

Using the oft-used Kennicutt-calibrated [O II]–SFR relation (Kennicutt 1992):

$$\text{SFR}([\text{O II}], M_{\odot} \text{ yr}^{-1}) = 1.4 \times 10^{-41} L_{\text{[O II]}} \text{(erg s}^{-1})$$

(1)
in the image plane, we find the SFR for COOL J1241+2219 to be $1700^{+800}_{-500} \ M_\odot \ yr^{-1}$. Instead, if we use the SFR–[O II] relation empirically corrected for stellar metallicity and takes as input the galaxy stellar mass (Equation (8) in Gilbank et al. 2010), we find the calculated SFR to be $4600^{+3400}_{-2400} \ M_\odot \ yr^{-1}$.

Without stronger constraints on metallicity and dust in our SED posterior distributions, their impact on the systematic uncertainties in the SFR–[O II] flux scaling relations cannot be mitigated (e.g., Kewley et al. 2004; Bicker & Fritz-v Alvensleben 2005). It is important to note that current constraints are within the metallicity values observed for local universe galaxies.

### 4.5. Lens Modeling and Magnification

We provide a brief summary of the gravitational lensing analysis used in this work here, and we refer the reader to Kneib & Soucail (1996), Richard et al. (2010), Verdugo et al. (2011) for a more in depth discussion of the lensing algorithm used. We adopt a parametric approach using LENSTOOL (Jullo et al. 2007) to model the mass distribution with dual pseudo-isothermal ellipsoids (dPIEs, Elíasdóttir et al. 2007), using an MCMC method to estimate the parameters and their uncertainties. The lens modeling strategy generally follows Sharon et al. (2020), with a cosmology with $\Omega_M = 0.3$, $\Omega_L = 0.7$. We model the lens plane with two dPIE potentials, one to represent a group-scale dark-matter halo, and one that represents the contribution from the main lens galaxy. The geometric parameters of the galaxy-scale halo are fixed to its observed properties (position, ellipticity, and position angle), while the slope and overall normalization are allowed to vary. The choice of two halos is motivated by the arc radius, which suggests the presence of a mass structure larger than a single massive galaxy, and the presence of other luminous early-type galaxies with colors consistent with the apparent central lens galaxy, again indicative of a more complex structure. Hence, the two dPIE clumps used here are combined to map the dark matter (DM) at the cluster scale and the central lens galaxy.

The smooth appearance of the arc in the ground-based data limits the availability of lensing constraints; a more complex lens plane, with additional galaxy-scale halos, is not required in order to satisfy these constraints. Hence, we do not include any other galaxy-scale halos in this initial model (see Figure 6). We choose to use the peak of the galaxy core as a constraint (see system 1 with four images in Figure 6), the surroundings of the star-forming clump merging on the critical curve (system 2), and a critical curve constraint (see green cross). The precision assumed for the location of the constraints in the lens model is $0.1^\prime$, with an rms of the optimized model in the image plane of $0.1^\prime.12$.

To estimate the magnification, we calculate the ratio of flux in the image plane using the GALFIT modeling image of the arc and its pixel-matched source plane reconstruction for 100 realizations of the lens model posterior probability distribution, resulting in an estimate of the magnification of $32^{+2}_{-5}$. We apply this as a division factor to the arc imaging in the image plane, to convert the derived physical parameter values to the source plane.

The uncertainties used in the SED fitting do not include the system-wide magnification uncertainty, which, when combined with the statistical uncertainties, yields a final stellar mass for COOL J1241+2219 of $\log M(M_\odot) = 10.11^{+0.21}_{-0.26}$. Similarly, the demagnified current star-formation rate (SFR, in units of $M_\odot \ yr^{-1}$) is calculated to be $27^{+3}_{-13}$ (see Figure 7), which makes this an active star-forming galaxy, given the constraints of the SED and lens modeling. This is assuming that the SFR in the age bin 0–50 Myr is a proxy for current star-forming activity in COOL J1241+2219.

#### Table 2

| Parameter | Description | Priors | Best-Fit Value |
|-----------|-------------|--------|----------------|
| log($M_{\text{star}}/M_\odot$) | Total stellar mass formed | Top hat: [10, 13] | $11.77^{+0.20}_{-0.25}$ |
| log(SFR$_{\text{adj}}$) | Ratio of the SFRs in adjacent age bins: nonparametric SFH | Student’s r: Bin 1: mean = 0.0, scale = 0.3, $\nu$ = 2 | $0.14^{+0.08}_{-0.05}$ |
| log($Z/M_\odot$) | Stellar metallicity in units of log($Z/M_\odot$) | Bin 2: mean = 0.0, scale = 0.3, $\nu$ = 2 | $0.10^{+0.01}_{-0.02}$ |
| $\tau_{3,2}$ | Diffuse dust optical depth | Top hat: [–2, 0.5] | $-1.32^{+0.06}_{-0.10}$ |
| spec$_\text{norm}$ | Factor by which to scale the spectrum to match photometry | Top hat: [0.01, 1.00] | $0.14^{+0.09}_{-0.13}$ |
| gas$_{\text{log}(Z/M_\odot)}$ | Gas-phase metallicity in units of log($Z/M_\odot$) | Top hat: [0.1, 3.0] | $1.94^{+0.50}_{-0.51}$ |
| log($U^\nu$) | Gas Ionization Parameter $U = n_\gamma/n_{H_2}$ | Top hat: [–4.0, –1.0] | $-0.20^{+0.03}_{-0.02}$ |

Note.

* Considered as nuisance parameters.
5. Discussion

5.1. Luminosity Function

We compare the brightness of COOL J1241+2219 to $M_{UV}^s$, the characteristic luminosity in the parametric luminosity function (Schechter 1976). As a proxy for the UV absolute magnitude, we compute $M_{1600}$, calculated from synthetic photometry applied to the best-fit SED model, using a top-hat filter from 1550 to 1650 Å. The UV continuum is roughly flat, and hence differences in definitions of UV magnitudes in the literature do not significantly impact the comparison of our reported absolute magnitude to published values.

Taking into account the lensing magnification, $M_{UV}$ for COOL J1241+2219 is $-22.2\pm0.2$. Bouwens et al. (2015) give a characteristic $M_{UV}^s$ at $z=5$ of $-21.1\pm0.2$ (see Table 6 in Bouwens et al. 2015; we use the higher uncertainty value). Hence, we find that for COOL J1241+2219, $M_{UV}^s-M_{UV}=1.0\pm0.5$. In other words, COOL J1241+2219 is approximately two to four times more luminous than an object with the characteristic UV luminosity at these redshifts.

5.2. The UV Continuum Slope

The observed UV continuum slope of star-forming galaxies is closely related to emission from massive stellar populations, as well as dust extinction in the galaxy. We fit the best-fit SED model for COOL J1241+2219 with a power law in $f_\lambda$, i.e., $f_\lambda \propto \lambda^{\beta}$, between rest-frame 1268 and 2580 Å (the entirety of the fitting windows described initially by Calzetti et al. 1994). For COOL J1241+2219, the best-fit UV continuum slope is calculated to be $\beta = -2.2 \pm 0.2$. This is on the blue side of the distribution of $\beta$ observed in (albeit fainter) galaxies at $z = 5$ in (e.g., Finkelstein et al. 2012) but consistent with their measured distribution. However, the mean relation between $\beta$ and UV luminosity (Bouwens et al. 2014) for galaxies at $z = 5$ predicts a value of $\beta$ for COOL J1241+2219 of $\sim-1.6$, significantly shallower than we observe. The observed $\beta$ is in fact consistent with that for galaxies $\sim$1 mag fainter, emphasizing that COOL J1241+2219 is unusually blue for its total UV luminosity.

5.3. Star-forming Main Sequence

The relationship between SFR and stellar mass of a galaxy is a well-studied relationship in galaxy populations and large samples, known as the star-forming main sequence (SFMS; Noeske et al. 2007; Rodighiero et al. 2011). It is informative to place these properties of COOL J1241+2219 in relation to the SFMS of a galaxy sample in the same epoch of observation. We compare the SFR and stellar mass of COOL J1241+2219 with galaxies at $4 < z < 6$ in the Hubble Space Telescope Frontier Fields. The SFMS for galaxies in this field in the redshift bins $4 < z < 5$, and $5 < z < 6$ has been reported in Santini et al. (2017). We direct the reader to the solid lines in bottom panels in Figure 2 of that publication, which reflect the best-fit SFMS, corrected for Eddington bias. The demagnified SFR and stellar mass of COOL J1241+2219 ($\log M_\odot$ = 10.11$^{+0.21}_{-0.26}$ and SFR (M_\odot yr$^{-1}$) = 27$^{+13}_{-9}$), are consistent with these power-law fits for the SFMS within 1σ, but incompatible with high SFRs corresponding to burst mode(s) of star formation during major mergers or in the densest regions. It is also important to note that the source plane properties of COOL J1241+2219 place this galaxy among the highest mass values reported in those data.

5.4. Stellar Populations in COOL J1241+2219

COOL J1241+2219 has a stellar mass consistent within a factor of a few of that of the present-day Milky Way (Licquia & Newman 2015), and is forming stars approximately one to two orders of magnitude faster. This rate of star formation is not surprising, given that the available time to build this galaxy at $z = 5.04$ is only $\sim$1 Gyr. COOL J1241+2219 is extremely UV bright with little apparent extinction. This stands in strong contrast to some similarly massive galaxies known from this early epoch (e.g., Riechers et al. 2013; Strandet et al. 2017; Casey et al. 2019), which tend to be much dustier with SFRs an order of magnitude larger. This difference is of course due to the selection process used to identify these objects. What is less clear, however, is whether COOL J1241+2219 looked like these dust-enshrouded examples of early star formation at any point in its history. The relative contribution of dust-enshrouded star formation to the total star-formation rate at this epoch is not well measured (Casey et al. 2019) and so objects like COOL J1241+2219 that are bright enough to allow detailed studies are of particular interest.

The current constraints from SED fitting (see Figures 5 and 7) imply that the SFH of COOL J1241+2219 is consistent with a constant SFR across the $\sim$1 Gyr available. The current data, however, are insufficient to say more; the nonparametric model forms the bulk of its stars in the oldest (and widest) bin, and a brief and much elevated star-formation episode in that interval is certainly possible.
Despite significant ongoing star formation, and a blue UV slope that suggests little extinction, COOL J1241+2219 shows no evidence for Ly$\alpha$ emission in the extracted spectrum. In addition, we carefully examine the stacked 2D spectrum, and see no evidence of spatially extended Ly$\alpha$ emission outside of the continuum. This complete lack of emission indicates the presence of sufficient dust to suppress Ly$\alpha$, as one would expect in a massive galaxy with a sustained SFH and resulting enrichment. These results are somewhat in tension. However, note that the measured scatter in galaxies between the Ly$\alpha$ escape fraction and continuum extinction is large (Hayes et al. 2011) and suppression of Ly$\alpha$ emission is possible even with extremely sparse dust columns.

6. Summary and Future Work

We report the discovery of COOL J1241+2219, a strongly lensed galaxy at $z = 5.043 \pm 0.002$ that is lensed by a moderate-mass galaxy cluster at $z = 1.00 \pm 0.001$. In this first work by the COOL-LAMPS collaboration, we characterize the lensed galaxy using $griz$H band photometry and optical and NIR spectroscopy. Using Prospector, we perform a stellar-population synthesis analysis spectrophotometrically, to constrain the stellar mass, SFH, stellar metallicity, and dust properties. Our dPIE-halo-based lens mass modeling from LENSTOOL, based on ground-based imaging, implies a median source magnification of at least $\approx 30$, which puts the stellar mass and star-formation rate (in the youngest age bin, closest to the epoch of observation) at $\log M_\star = 10.11^{+0.21}_{-0.26}$ and SFR = $27^{+3}_{-9} M_\odot yr^{-1}$, respectively. This places this galaxy on the star-forming main sequence, at the massive end of that relation, and we find that COOL J1241+2219 is two to four times more luminous than a galaxy with the characteristic UV luminosity at these redshifts. We constrain a star-formation history for COOL J1241+2219 starting at $z \approx 20$ that is consistent with constant star formation across $\approx 1$ Gyr of cosmic time.

Current data are insufficient to constrain the metallicity and star-formation history robustly. Also, as is almost always true, the lens model built on ground-based data limits the precision and accuracy of some aspects of our analysis; a lens model built from much sharper imaging (e.g., with HST) will provide insights into fluxes, flux variations, and identification of a variety of faint spatial features within this galaxy much more comprehensively via a proper source plane reconstruction. Moreover, our current spectroscopy only samples the apparently brightest clump in this galaxy. Once lensed into the observed giant arc, COOL J1241+2219 is the brightest galaxy at $z > 5$ currently known, and thus presents a unique opportunity for the detailed study of a massive galaxy immediately after the epoch of reionization. COOL J1241+2219 is bright enough that precise high resolution observations will be possible across a broad range of wavelengths.

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ORCID iDs

Guillaume Mahler @ https://orcid.org/0000-0003-3266-2001
Jane R. Rigby @ https://orcid.org/0000-0002-7627-6551
Keren Sharon @ https://orcid.org/0000-0002-7559-0864
Antony A. Stark @ https://orcid.org/0000-0002-2718-9996

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