Minimization variational principles for acoustics, elastodynamics and electromagnetism in lossy inhomogeneous bodies at fixed frequency

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The classical energy minimization principles of Dirichlet and Thompson are extended as minimization principles to acoustics, elastodynamics and electromagnetism in lossy inhomogeneous bodies at a fixed frequency. This is done by building upon the ideas of Cherkaev and Gibiansky, who derived minimization variational principles for quasistatics. In the absence of free current, the primary electromagnetic minimization variational principles have a minimum, which is the time-averaged electrical power dissipated in the body. The variational principles provide constraints on the boundary values of the fields when the moduli are known. Conversely, when the boundary values of the fields have been measured, then they provide information about the values of the moduli within the body. This should have applications to electromagnetic tomography. We also derive saddle-point variational principles that correspond to the variational principles of Gurtin, Willis and Borcea.

Keywords: variational principles; acoustics; elastodynamics; electromagnetism

1. Introduction

As our goal is to extend the classical minimization principles to acoustics, elastodynamics and electromagnetism, let us begin by recalling the well-known classical results of electrostatics. Consider the equation of electrostatics

\[ \nabla \cdot \varepsilon \nabla V = 0 \]  

for a real-valued potential \( V(\mathbf{x}) \) in a body \( \Gamma \) containing a locally isotropic material with scalar real positive permittivity \( \varepsilon(\mathbf{x}) \). This equation can be rewritten in the equivalent form

\[ D = \varepsilon E, \quad E = -\nabla V \quad \text{and} \quad \nabla \cdot D = 0, \]  

where the first relation is known as the constitutive relation linking the electric displacement field \( D(\mathbf{x}) \) and the electrical field \( E(\mathbf{x}) \), where these two fields

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satisfy the differential constraints embodied in the second two relations. Two variational principles are well known. We have the Dirichlet variational principle that

\[ W(V) = \inf_V W(V), \quad \text{where } W(V) = \int_{\Omega} \epsilon |\nabla V|^2/2, \]

(1.3)

in which the infimum is taken over all trial potentials \( V(x) \), satisfying the boundary condition that \( V(x) = V_0(x) \) on the boundary \( \partial \Omega \). Here, and subsequently, it is implicitly assumed that integrals over \( \Omega \) are volume integrals \( dv \) and, similarly, it will be implicitly assumed that integrals over \( \partial \Omega \) are surface integrals \( ds \). We also have a dual Thompson variational principle that

\[ \tilde{W}(D) = \inf_D \tilde{W}(D), \quad \text{where } \tilde{W}(D) = \int_{\Omega} |D|^2/(2\epsilon), \]

(1.4)

where the infimum is taken over all divergence-free trial displacement fields \( D \), satisfying the boundary condition that \( D \cdot n = D \cdot n \) on \( \partial \Omega \), in which \( n \) is the outward unit normal to the surface \( \partial \Omega \). The quantity \( \epsilon |\nabla V|^2/2 = |D|^2/(2\epsilon) \) is the electrostatic energy density inside the body, and so these variational principles are also known as energy minimization principles. In the mathematically analogous DC electrical conductivity problem, where \( \epsilon(x) \) and \( D(x) \) are replaced by the conductivity \( \sigma(x) \) and the current \( J(x) \), this quantity is half the electrical power dissipation and, in that context, the variational principles could more accurately be called power dissipation minimization principles.

An appealing feature of these variational principles is that the minimum values \( W(V) \) and \( \tilde{W}(D) \) can be expressed in terms of the values of \( V(x) \) and \( D \cdot n \) at the boundary \( \partial \Omega \),

\[ W(V) = \tilde{W}(D) = -\frac{1}{2} \int_{\partial \Omega} (D \cdot n) V(x). \]

(1.5)

So these variational principles provide a constraint on the values of \( V_0(x) = V(x) \) and \( q_0(x) = -D(x) \cdot n \) at the boundary \( \partial \Omega \). The relationship between these boundary functions \( q_0 \) and \( V_0 \) is linear and the relationship \( q_0 = N V_0 \) defines the Dirichlet to Neumann map \( N \), which is a real self-adjoint positive semi-definite operator. Therefore, \( N \) is determined by knowledge of the quadratic form

\[ w(V_0) = \langle V_0, N V_0 \rangle/2, \]

(1.6)

where, for any two scalar or vector, real or complex, valued functions \( p(x) \) and \( q(x) \) defined on the boundary \( \partial \Omega \), we denote

\[ \langle p, q \rangle = \int_{\partial \Omega} p(x) \cdot q(x). \]

(1.7)

Given trial potentials \( V(x, V_0) \), parametrized by the function \( V_0 \), with \( V(x, V_0) = V_0(x) \) on \( \partial \Omega \), the Dirichlet variational principle implies the inequality

\[ w(V_0) \leq W(V), \]

(1.8)

and provides an ‘upper bound’ on the map \( N \). By upper bound we mean an upper bound on the associated non-negative quadratic form \( w(V_0) \).
To obtain a ‘lower bound’, one uses the Thompson variational principle. Since \( w(V_0) = 0 \) when \( V_0 \) is constant, the null space of \( N \) consists of all constant fields, and so \( N \) is determined by knowledge of the quadratic form

\[
\bar{w}(q_0) = \langle q_0, N^{-1} q_0 \rangle / 2,
\]

for all fields \( q_0(x) \) with zero average value on \( \partial \Omega \). The functionals \( w(V_0) \) and \( \bar{w}(q_0) \) are Legendre transforms of each other,

\[
w(V_0) = \inf_{q} \{ \langle V_0, q \rangle - \bar{w}(q) \},
\]

where the infimum is over all fields \( q(x) \) with zero average value on \( \partial \Omega \), and so an upper bound on the functional \( \bar{w}(q) \) provides a lower bound on \( w(V_0) \). Given trial displacement fields \( D(x, q_0) \) parametrized by the function \( q_0 \), with \( -D(x) \cdot n = q_0(x) \) on \( \partial \Omega \), the Thompson variational principle implies the inequality

\[
\bar{w}(q_0) \leq \bar{W}(D),
\]

and provides a lower bound on the map \( N \).

Thus, the Dirichlet and Thompson variational principles allow one to bound the Dirichlet to Neumann map \( N \). As shown by Berryman & Kohn (1990), they can also be used in an inverse fashion: if \( N \) and hence the functionals \( w(V_0) \) and \( \bar{w}(q_0) \) are known, one can obtain information about the permittivity \( \varepsilon(x) \) from the inequalities (1.8) and (1.11) by appropriately choosing the trial potentials \( V(x, V_0) \) and trial displacement fields \( D(x, q_0) \). It has been proved that knowledge of \( N \) uniquely determines \( \varepsilon(x) \) when \( \varepsilon(x) \) is piecewise analytic (Kohn & Vogelius 1984), or more generally when \( \varepsilon(x) \) is smooth (Sylvester & Uhlmann 1987).

In this paper, we extend these variational principles to acoustics, elastodynamics and electromagnetism in lossy inhomogeneous bodies at fixed frequency. By lossy we mean that the time-averaged dissipation of mechanical or electrical energy into heat is positive everywhere, which means that the imaginary part of certain tensors, such as the permittivity tensor, are positive definite. As with the Dirichlet and Thompson variational principles, the variational principles we obtain are minimization (not saddle-point or stationary) variational principles. They apply even if the moduli entering the constitutive equations are discontinuous across some interior interfaces. Here, we explicitly treat only the case where the moduli are continuous since the discontinuous case can be obtained by taking limits.

The variational principles we derive in this paper correspond to minimizing convex functionals. Hence, they give a standard and reliable mathematical framework: existence and uniqueness of their solution can be stated in an appropriate functional space (depending on whether one deals with acoustics, elastodynamics or electromagnetism, one is led to different functional spaces together with the various associated boundary conditions). Furthermore, estimates for the solution can be easily established and limit analysis of related perturbation problems can be handled using the tools of variational analysis, such as Gamma-convergence (Braides 2002). However, we do not consider such questions here.

A great variety of stationary variational principles for acoustics, elastodynamics and electromagnetism have been derived, some allowing for coupling between the various fields and most not assuming time harmonicity (e.g. Oden & Reddy (1983).
and Altay & Dökmeci (2004) and references therein). However, to our knowledge, none of these are minimizing variational principles at fixed frequency. Among the stationary variational principles, we draw particular attention to the variational principles of Gurtin (1964) and Willis (1981, 1984; see also Tao 1966). These are minimizing variational principles in the Laplace domain, but not in the frequency domain. We recover the Gurtin–Willis variational principles in the time-harmonic setting where they correspond to saddle-point variational principles. We remark that Ben-Amoz (1966) and Willis (1981, 1984) derived variational principles that are dynamic analogues of the Hashin–Shtrikman (1962a–c, 1963) variational principles, involving trial polarization fields. We do not consider such variational principles here. Of course, given a (inhomogeneous) linear differential operator $\mathcal{A}$, mapping $m$-component complex vector fields to $m$-component complex vector fields, there is a trivial minimization variational principle associated with the equation $\mathcal{A}u = h$ for the $m$-component potential $u$, in which $h$ is an $m$-component source term. When the equation $\mathcal{A}u = h$, with $u = u_0$ on $\partial T$, has a unique solution for $u$, the infimum in

$$0 = \inf_u \int_T (\mathcal{A}u - h) \cdot (\mathcal{A}u - h),$$

where the overline denotes complex conjugation and the infimum is over all fields $u$, with $u = u_0$ on $\partial T$, is clearly attained when $u = u$. What separates such variational principles from those of Dirichlet and Thompson is that, as explained above, the latter variational principles give us useful information on the boundary values of the fields. We desire variational principles with this property. Also, from a mathematical viewpoint, the existence of a solution cannot be deduced from (1.12), and in case there is a solution no interesting inequality results from (1.12).

In an important development, Cherkaev & Gibiansky (1994) extended the Dirichlet and Thompson variational principles to the quasistatic electromagnetic equations (and also to the quasistatic elastodynamic equations), where equations (1.1) and (1.2) still hold, but $\varepsilon(x)$, $V(x)$, $E(x)$ and $D(x)$ are complex and the imaginary part of $\varepsilon(x)$ is positive. A measure of the success of their variational principles is that they led to tight bounds on the complex effective permittivity (Milton 1990) and effective viscoelastic moduli (Gibiansky & Milton 1993; Milton & Berryman 1997; Gibiansky et al. 1999) of composites. Since we build upon the work of Cherkaev and Gibiansky, let us briefly review their ideas in this context.

By rewriting the constitutive law $D = \varepsilon E$ in terms of its real and imaginary parts

$$\begin{pmatrix} D'' \\ D' \end{pmatrix} = \begin{pmatrix} \varepsilon'' & \varepsilon' \\ \varepsilon' & -\varepsilon'' \end{pmatrix} \begin{pmatrix} E' \\ E'' \end{pmatrix},$$

where the prime denotes the real part and the double prime denotes the imaginary part, it is evident that, since $\varepsilon''(x) > 0$, one has the Cherkaev–Gibiansky saddle-point variational principle

$$Q(V', V'') = \inf_{V'} \sup_{V''} Q(V', V'') = \sup_{V'} \inf_{V''} Q(V', V''),$$

where

$$Q(V', V'') = \int_T \begin{pmatrix} \nabla V' \\ \nabla V'' \end{pmatrix} \cdot \begin{pmatrix} \varepsilon'' & \varepsilon' \\ \varepsilon' & -\varepsilon'' \end{pmatrix} \begin{pmatrix} \nabla V' \\ \nabla V'' \end{pmatrix},$$

Proc. R. Soc. A (2009)
and the infimum and supremum are over all trial potentials, with $V'(x) = V'(x)$ and $V''(x) = V''(x)$ on the boundary $\partial \Gamma$. To obtain a minimization variational principle from this saddle-point variational principle, Cherkaev & Gibiansky (1994) made a partial Legendre transform of the saddle-shaped quadratic form to convert it into a convex quadratic form. This is equivalent to rewriting the constitutive law in the form

$$
\begin{pmatrix}
  D'' \\
  E''
\end{pmatrix}
= \mathcal{E}
\begin{pmatrix}
  E' \\
  -D'
\end{pmatrix},
$$

where the matrix

$$
\mathcal{E} =
\begin{pmatrix}
  \varepsilon''' + (\varepsilon')^2 / (\varepsilon'') & \varepsilon' / \varepsilon'' \\
  \varepsilon' / \varepsilon'' & 1 / \varepsilon''
\end{pmatrix}
$$

is positive definite when $\varepsilon''(x)$ is positive. From this, it is evident that one has the Cherkaev–Gibiansky minimization variational principle

$$
Y(V', D') = \inf_{V'} \inf_{D'} Y(V', D'),
$$

where

$$
Y(V', D') = \int_{\Omega} \left( \nabla V \right) \cdot \mathcal{E} \left( \nabla V \right) - \varepsilon \cdot D',
$$

and the infimums are over all trial real potentials, with $V'(x) = V'(x)$ on $\partial \Gamma$, and all trial real divergence-free displacement fields, with $D \cdot n = D \cdot n$ on $\partial \Gamma$.

As an aside, we remark that, as shown in §18 of Milton (1990), one can extend these ideas of Gibiansky and Cherkaev to obtain variational principles for problems where the constitutive law takes the form $D = \varepsilon E$, where the matrix-valued field $\varepsilon(x)$ is not symmetric. For example, for conduction in a fixed magnetic field, the conductivity tensor $\sigma(x)$ is real but not symmetric. Expressing $\sigma(x) = \sigma_s(x) + \sigma_a(x)$, where $\sigma_s(x) = \sigma_s^T(x)$ and $\sigma_a(x) = -\sigma_a^T(x)$ are the symmetric and antisymmetric parts of the conductivity tensor, respectively, one considers current fields $J(x)$ and electric fields $E(x)$ that solve the conductivity equations,

$$
J = (\sigma_s + \sigma_a) E, \quad E = -\nabla V \quad \text{and} \quad \nabla \cdot J = 0,
$$

in conjunction with current fields $\hat{J}(x)$ and electric fields $\hat{E}(x)$ that solve the adjoint problem,

$$
\hat{J} = (\sigma_s - \sigma_a) \hat{E}, \quad \hat{E} = -\nabla \hat{V} \quad \text{and} \quad \nabla \cdot \hat{J} = 0.
$$

By adding and subtracting the two constitutive laws, one obtains a new constitutive law

$$
\begin{pmatrix}
  J_s \\
  -J_a
\end{pmatrix}
= \begin{pmatrix}
  \sigma_s & \sigma_a \\
  -\sigma_a & -\sigma_s
\end{pmatrix}
\begin{pmatrix}
  E_s \\
  E_a
\end{pmatrix},
$$

involving fields

$$
J_s \equiv J + \hat{J}, \quad J_a \equiv J - \hat{J}, \quad E_s \equiv E + \hat{E} \quad \text{and} \quad E_a \equiv E - \hat{E},
$$
which satisfy the differential constraints

\[ \nabla \cdot J_s = \nabla \cdot J_a = 0, \quad E_s = -\nabla V_s \quad \text{and} \quad E_a = -\nabla V_a, \quad (1.24) \]

where \( V_s \equiv V + \hat{V} \) and \( V_a \equiv V - \hat{V} \). Conversely, given any fields that solve (1.22) and (1.24), one can recover the associated fields that solve (1.20) and (1.21). The two problems are equivalent. Since the matrix entering the constitutive law (1.22) is real and symmetric, there is clearly an associated saddle-point variational principle when \( \sigma_s(x) \) is positive definite for all \( x \). Fannjiang & Papanicolaou (1994) showed how this variational principle provides useful asymptotic estimates for the mathematically analogous problem of convection-enhanced diffusion. Also, by rewriting the constitutive law (1.22) in the equivalent form

\[ (J_s) = \begin{pmatrix} \sigma_s - \sigma_a \sigma_s^{-1} \sigma_a & \sigma_a \sigma_s^{-1} \\ -\sigma_s^{-1} \sigma_a & \sigma_s^{-1} \end{pmatrix} \begin{pmatrix} E_s \\ J_a \end{pmatrix}, \quad (1.25) \]

which involves a matrix that is symmetric and positive definite when \( \sigma_s(x) \) is positive definite for all \( x \), one obtains an associated minimization variational principle (Milton 1990). In this paper, we do not treat such cases where the constitutive laws involve non-symmetric matrices; this will be covered elsewhere.

2. General theory

Let \( d \) be the dimension of the space and let \( \mathcal{F}(x) \) and \( \mathcal{G}(x) \) be two complex-valued fields of the form

\[ \mathcal{F} = \begin{pmatrix} F \\ f \end{pmatrix} \quad \text{and} \quad \mathcal{G} = \begin{pmatrix} G \\ g \end{pmatrix}, \quad (2.1) \]

where \( F(x) \) and \( G(x) \) are \( m \times d \) dimensional matrix fields and \( f(x) \) and \( g(x) \) are \( m \)-component vector fields, and which satisfy the differential constraints

\[ \mathcal{F} = \nabla u \quad \text{and} \quad h + \nabla \mathcal{G} = 0, \quad (2.2) \]

where \( u(x) \) is an \( m \)-component potential and \( h(x) \) is some \( m \)-component source term, and

\[ \nabla u \equiv \begin{pmatrix} \nabla u \\ u \end{pmatrix} \quad \text{and} \quad \nabla \mathcal{G} \equiv -\nabla \cdot \mathcal{G} + g \quad (2.3) \]

serve to define the operators \( \nabla \) and \( \nabla \cdot \). The differential constraints imply the key property that, on any domain \( \Gamma \),

\[ \int_{\Gamma} \mathcal{G} \cdot \mathcal{F} + h \cdot u = \int_{\partial \Gamma} (G \cdot n) \cdot u, \quad (2.4) \]

as follows from integration by parts, where

\[ \mathcal{G} \cdot \mathcal{F} \equiv G \cdot F + g \cdot f \quad \text{and} \quad G \cdot F \equiv \text{Tr}(G^T F). \quad (2.5) \]

(Note that (2.5) is an inner product only when \( \mathcal{G} \) and \( \mathcal{F} \) are real, since it does not involve complex conjugation.) Here, following the standard notation, we define

\[ \{\nabla u\}_{ij} = \frac{\partial u_i}{\partial x_j} \quad \text{and} \quad \{\nabla \cdot G\}_i = \sum_{j=1}^{d} \frac{\partial G_i}{\partial x_j}. \quad (2.6) \]
Now suppose the fields $\mathcal{F}$ and $\mathcal{G}$ are linked by the constitutive relation
\[ \mathcal{G}(x) = Z(x)\mathcal{F}(x), \tag{2.7} \]
where $Z(x) = Z'(x) + iZ''(x)$ is a complex symmetric linear map (which is represented by a complex symmetric $m(d+1) \times m(d+1)$ matrix when $\mathcal{G}$ and $\mathcal{F}$ are represented by $m(d+1)$-component vectors). The map $Z$ can be expressed in the block form as
\[ Z = \begin{pmatrix} L & K \\ K^T & M \end{pmatrix}, \tag{2.8} \]
where (for each $x$) $L(x)$ is a symmetric linear map on the space of $m \times d$-dimensional matrices; $K(x)$ is a linear map from $m$-component vectors to $m \times d$-dimensional matrices; $K^T(x)$ is a linear map from $m \times d$-dimensional matrices to $m$-component vectors; and $M(x)$ is a symmetric linear map on the space of $m$-component vectors. Thus, written in terms of components, (2.7) implies
\[ G_{ij} = L_{ijk}F_{k\ell} + K_{ijk}f_{k\ell} \quad \text{and} \quad g_i = K_{ik}F_{k\ell} + M_{ik}f_{k\ell}, \tag{2.9} \]
in which sums of repeated indices are assumed. Since $L$ and $M$ are symmetric maps, we have
\[ L_{ijk\ell} = L_{k\ell ij} \quad \text{and} \quad M_{ik} = M_{ki}. \tag{2.10} \]

The constitutive law and differential constraints imply that $u$ satisfies
\[ h + \nabla(Z \nabla u) = 0, \tag{2.11} \]
or, equivalently,
\[ \nabla \cdot (L\nabla u + Ku) = h + K^T\nabla u + Mu. \tag{2.12} \]
Our analysis thus applies to equations of this form, with $L$ and $M$ having the symmetries (2.10).

Let $(u, \mathcal{F}, \mathcal{G})$ be a solution of (2.2) and (2.7). Taking real and imaginary parts, the differential constraints imply
\[ F' = \nabla u', \quad h' + \nabla G' = 0, \quad F'' = \nabla u'' \quad \text{and} \quad h'' + \nabla G'' = 0, \tag{2.13} \]
where the prime denotes the real part and the double prime the imaginary part, and the constitutive law can be written in the form
\[ \begin{pmatrix} \mathcal{G}'' \\ \mathcal{G}' \end{pmatrix} = \begin{pmatrix} Z'' & Z' \\ Z' & -Z'' \end{pmatrix} \begin{pmatrix} \mathcal{F}' \\ \mathcal{F}'' \end{pmatrix}. \tag{2.14} \]

To make use of this representation, let us further suppose that the imaginary part of $Z$ satisfies
\[ Z''(x) \geq \alpha I, \quad \forall x \in \Gamma, \tag{2.15} \]
for some $\alpha > 0$. Extending the ideas of Cherkaev & Gibiansky (1994) consider the real-valued quadratic functional
\[ Q(u', u'') = \int_{\Gamma} \begin{pmatrix} \nabla u' \\ \nabla u'' \end{pmatrix} \cdot \begin{pmatrix} Z'' & Z' \\ Z' & -Z'' \end{pmatrix} \begin{pmatrix} \nabla u' \\ \nabla u'' \end{pmatrix} + 2(h'' \cdot u' + h' \cdot u'') \]
\[ = \int_{\Gamma} (\nabla u) \cdot Z \nabla u + 2h \cdot u'', \tag{2.16} \]

Proc. R. Soc. A (2009)
subject to the boundary conditions that \( u' = u' \) and \( u'' = u'' \) on \( \partial \Gamma \). Let \( s(x) \) be a real-valued \( m \)-component vector field, with \( s = 0 \) on \( \partial \Gamma \). Now \( Q(u' + s, u'') \) is a convex quadratic function of \( s \) (with a linear component), while \( Q(u', u' + s) \) is a concave quadratic function of \( s \) (with a linear component). Therefore, on the space of fields satisfying the boundary conditions, the pair \((u', u'')\) that is at the saddle point of the variational problem

\[
\inf \sup Q(u', u'')
\]  

exists and is unique and corresponds to the stationary point of \( Q(u', u'') \), and is also at the saddle point of the variational problem

\[
\sup \inf Q(u', u'').
\]

Now let us prove we have the variational principle

\[
Q(u', u'') = \inf \sup Q(u', u'') = \sup \inf Q(u', u''),
\]  

where the infimum and supremum are over fields with \( u' \) and \( u'' \), with any fixed values \( u'_0 \) and \( u''_0 \) at the boundary \( \partial \Gamma \), and \( u' \) and \( u'' \) are the solutions with these prescribed boundary conditions.

Let \((u, F, G)\) be a solution of (2.2) and (2.7). As \( s = 0 \) on \( \partial \Gamma \), in a similar way as (2.4) was derived, we obtain

\[
\int_{\Gamma} G'' \cdot (\nabla s) + h'' \cdot s = 0 \quad \text{and} \quad \int_{\Gamma} G' \cdot (\nabla s) + h' \cdot s = 0.
\]

Thus, we note that

\[
Q(u' + s, u'') = Q(u', u'') + \int_{\Gamma} 2G'' \cdot (\nabla s) + 2h'' \cdot s + (\nabla s) \cdot Z'' \nabla s
\]

\[
= Q(u', u'') + \int_{\Gamma} (\nabla s) \cdot Z'' \nabla s,
\]

and, in the same way, we obtain

\[
Q(u', u'' + s) = Q(u', u'') - \int_{\Gamma} (\nabla s) \cdot Z'' \nabla s.
\]

So we note that the saddle point of (2.17) is at \((u', u'') = (u', u'')\).

Conversely, let \((u', u'')\) be a saddle point of (2.17), satisfying some boundary conditions \( u' = u'_0 \) and \( u'' = u''_0 \) on \( \partial \Gamma \). Consider

\[
\inf \sup Q(u' + s, u'' + t),
\]

where the infimum and supremum are taken over fields \( s \) and \( t \), satisfying \( s = t = 0 \) on \( \partial \Gamma \). A necessary condition for the saddle point to be at \( s = t = 0 \) is that the first-order derivative of the functional \( Q \) vanishes, i.e.

\[
\int_{\Gamma} 2 \left( \nabla (\nabla s) \right) \cdot \left( \frac{G''}{G'} \right) + 2h'' \cdot s + 2h' \cdot t = 0,
\]

\(\text{Proc. R. Soc. A (2009)}\)
for any $s$ and $t$, satisfying $s = t = 0$ on $\partial \Gamma$, where we have introduced
\[
\begin{pmatrix}
G''
\end{pmatrix}
= \begin{pmatrix}
Z'' & Z' \\
Z' & -Z''
\end{pmatrix}
\begin{pmatrix}
\nabla u'
\end{pmatrix}.
\tag{2.25}
\]

Upon setting
\[
G' = \begin{pmatrix}
G' \\
g'
\end{pmatrix}	ext{ and } G'' = \begin{pmatrix}
G''
\end{pmatrix},
\tag{2.26}
\]
the condition (2.24) reduces to
\[
\int_{\Gamma} G'' : \nabla s + (g'' + h'') \cdot s + G' : \nabla t + (g' + h') \cdot t = 0.
\tag{2.27}
\]

Integrating this by parts gives
\[
\int_{\Gamma} (-\nabla \cdot G'' + g'' + h'') \cdot s + (-\nabla \cdot G' + g' + h') \cdot t = 0,
\tag{2.28}
\]
and this will be satisfied for any $s$ and $t$, satisfying $s = t = 0$ on $\partial \Gamma$ if and only if
\[
h' + \nabla G' = 0 \text{ and } h'' + \nabla G'' = 0.
\tag{2.29}
\]

Thus, the Euler–Lagrange equations associated with the saddle-point variational principle coincide with the original equations (2.2).

There is also a dual saddle-point variational principle, analogous to the Thompson variational principle. By taking real and imaginary parts of the constitutive relation $F = KG$, where $K = \mathbf{Z}^{-1} = \mathbf{Z}^{-\dagger}$ has a negative definite imaginary part, we obtain
\[
\begin{pmatrix}
F'' \\
F'
\end{pmatrix} = \begin{pmatrix}
K'' & K' \\
K' & -K''
\end{pmatrix}
\begin{pmatrix}
G' \\
g'
\end{pmatrix}.
\tag{2.30}
\]

The dual variational principle involves the quadratic form
\[
R(G', G'') = \int_{\Gamma} \begin{pmatrix}
G' \\
g'
\end{pmatrix} \cdot \begin{pmatrix}
K'' & K' \\
K' & -K''
\end{pmatrix}
\begin{pmatrix}
G' \\
g'
\end{pmatrix} = \int_{\Gamma} (G \cdot KG)',
\tag{2.31}
\]
where $G'$ and $G''$ are given by
\[
G' = \begin{pmatrix}
\nabla \cdot G' - h'
\end{pmatrix} \text{ and } G'' = \begin{pmatrix}
\nabla \cdot G'' - h''
\end{pmatrix},
\tag{2.32}
\]
and $G'$ and $G''$ satisfy the boundary conditions
\[
G' \cdot n = G' \cdot n \text{ and } G'' \cdot n = G'' \cdot n \text{ on } \partial \Gamma,
\tag{2.33}
\]
where $G'$ and $G''$ are the real and imaginary parts of $G$. Note that the form of $G'$ given by (2.32) ensures that $h + \nabla G = 0$ inside $\Gamma$. By similar argument to the one establishing (2.19), we have the dual variational principle
\[
R(G', G'') = \sup_{G'} \inf_{G''} R(G', G'') = \inf_{G''} \sup_{G'} R(G', G''),
\tag{2.34}
\]
where the supremum and infimum are over fields satisfying (2.33).
Again, building upon the ideas of Cherkaev & Gibiansky (1994), let us rewrite the constitutive relation in the form

\[
\begin{pmatrix}
G'' \\
F''
\end{pmatrix} = \mathcal{L} \begin{pmatrix}
F' \\
-\mathcal{G}'
\end{pmatrix},
\]

where straightforward algebra shows that

\[
\mathcal{L} = \begin{pmatrix}
Z'' + Z'(Z'')^{-1}Z' & Z'(Z'')^{-1} \\
(Z'')^{-1}Z' & (Z'')^{-1}
\end{pmatrix}.
\]

(2.35)

For all \(x \in \Gamma\), the matrix \(\mathcal{L}(x)\) is positive definite since, using (2.14) and (2.35), we note that the associated quadratic form

\[
(F' \cdot G'' - F'' \cdot \mathcal{G}') = F' \cdot Z'' F' + F'' \cdot Z'' F''
\]

is non-negative, and zero only when \(F' = F'' = 0\), i.e. when \(F' = \mathcal{G}' = 0\).

Now consider the convex functional

\[
Y(u', G') = \int_{\Gamma} \left( \nabla u' \right) \cdot \mathcal{L} \left( \begin{pmatrix} \nabla u' \\ -\mathcal{G}' \end{pmatrix} \right) + 2h'' \cdot u',
\]

(2.36)

where \(\mathcal{G}'\) is given by

\[
\mathcal{G}' = \begin{pmatrix} G' \\ \nabla \cdot G' - h' \end{pmatrix}
\]

(to ensure that \(h' + \nabla \mathcal{G}' = 0\) inside \(\Gamma\)), and \(u'\) and \(G'\) are real and satisfy the boundary conditions that

\[
u' = u'_0 \quad \text{and} \quad G' \cdot n = G'_0 \cdot n \quad \text{on} \ \partial \Gamma.
\]

(2.40)

Let us prove we have the variational principle

\[
Y(u', G') = \inf_{u'} \inf_{G'} Y(u', G'),
\]

(2.41)

where the infimums are over all trial fields \(u'\) and \(G'\), satisfying the boundary conditions (2.40). Let \((u, G)\) be a solution of (2.2) and (2.7), with \(u' = u'_0\) and \(G' \cdot n = G'_0 \cdot n\) on \(\partial \Gamma\). Suppose we are given any \(m\)-component vector field \(s(x)\), with \(s = 0\) on \(\partial \Gamma\), and also an \(m \times d\) matrix-valued field \(T(x)\), with \(T \cdot n = 0\) on \(\partial \Gamma\). Then, (2.20) and integration by parts imply that

\[
\int_{\Gamma} 2G'' \cdot (\nabla s) + 2h'' \cdot s - 2T \cdot (\nabla u'') = 0, \quad \text{where} \ T = \begin{pmatrix} T \\ \nabla \cdot T \end{pmatrix}.
\]

(2.42)

From this identity, it follows that

\[
Y(u' + s, G' + T) = Y(u', G') + \int_{\Gamma} \left( \nabla s \right) \cdot \mathcal{L} \left( \begin{pmatrix} \nabla s \\ -T \end{pmatrix} \right),
\]

(2.43)

which implies \((u', G')\) is a minimizer of the variational principle.
Conversely, let \((u', G')\) be a minimizer of the variational principle, satisfying some boundary conditions \(u' = u_0'\) and \(G' \cdot n = G_0' \cdot n\) on \(\partial \Gamma\). Consider

\[
\inf_s \inf_T Y(u' + s, G' + T),
\]

where the infimum is over all fields \(s\) and \(T\), satisfying \(s = 0\) and \(T \cdot n = 0\) on \(\partial \Gamma\). A necessary condition for the infimum to be at \(s = T = 0\) is that the first-order derivative of the functional \(Y\) vanishes, i.e.

\[
\int_I 2 \left( \nabla s - T \right) \cdot \left( \nabla'' + 2 h'' \cdot s = 0, \tag{2.45} \right.
\]

for any \(s\) and \(T\), satisfying \(s = 0\) and \(T \cdot n = 0\) on \(\partial \Gamma\), where we have introduced

\[
T \equiv \left( \begin{array}{c} T \\ \nabla \cdot T \end{array} \right) \quad \text{and} \quad \left( \begin{array}{c} G'' \\ \mathcal{F}'' \end{array} \right) \equiv \mathcal{L} \left( \begin{array}{c} \nabla u' \\ -G' \end{array} \right). \tag{2.46}
\]

Upon setting

\[
\mathcal{F}'' = \left( \begin{array}{c} F'' \\ f'' \end{array} \right) \quad \text{and} \quad \mathcal{G}'' = \left( \begin{array}{c} G'' \\ g'' \end{array} \right), \tag{2.47}
\]

the condition (2.45) reduces to

\[
\int_I G'' \cdot \nabla s + (g'' + h'') \cdot s - T \cdot F'' - (\nabla \cdot T) \cdot f'' = 0. \tag{2.48}
\]

Integrating this by parts gives

\[
\int_I (\nabla \cdot G'' + g'' + h'') \cdot s - T \cdot (F'' - \nabla f'') = 0, \tag{2.49}
\]

and this will be satisfied for any \(s\) and \(T\), satisfying \(s = 0\) and \(T \cdot n = 0\) on \(\partial \Gamma\), if and only if

\[
h'' + \nabla \mathcal{G}'' = 0 \quad \text{and} \quad \mathcal{F}'' = \nabla \mathcal{F}'' . \tag{2.50}
\]

Upon identifying \(f''\) with \(u''\), we note that the Euler–Lagrange equations associated with the minimization variational principle coincide with the original equations (2.2).

The convex variational principle we have just stated is, as usual, associated with a dual variational principle

\[
\tilde{Y}(u'', G'') = \inf_{u''} \inf_{G''} \tilde{Y}(u'', G''), \tag{2.51}
\]

where

\[
\tilde{Y}(u'', G'') = \int_I \left( \begin{array}{c} G'' \\ \nabla \cdot G'' - h'' \end{array} \right) \cdot \mathcal{L}^{-1} \left( \begin{array}{c} G'' \\ \nabla u'' \end{array} \right) - 2 h' \cdot u'', \tag{2.52}
\]

in which

\[
\mathcal{G}'' = \left( \begin{array}{c} G'' \\ \nabla \cdot G'' - h'' \end{array} \right), \tag{2.53}
\]
and the infimum in (2.51) is over real fields \(u''\) and \(G''\), with \(u'' = u'\) and \(G'' \cdot n = G'' \cdot n\) on \(\partial \Gamma\). Since

\[
\mathcal{L}^{-1} = \begin{pmatrix}
(Z'')^{-1} & -(Z'')^{-1}Z' \\
-Z'(Z'')^{-1} & Z'' + Z'(Z'')^{-1}Z'
\end{pmatrix}, \tag{2.54}
\]

the expression for \(\tilde{Y}(u'', G'')\) can be written equivalently as

\[
\tilde{Y}(u'', G'') = \int_{\Gamma} \left( \frac{\nabla u''}{-G''} \right) \cdot \mathcal{L} \left( \frac{\nabla u''}{-G''} \right) - 2h' \cdot u''. \tag{2.55}
\]

These two minimization variational principles are not the only ones that can be derived. Indeed, there is a continuous two-parameter family of related variational principles. The differential constraints (2.2) and the constitutive relation (2.7) imply

\[
\tilde{F} = \nabla \tilde{u}, \quad \tilde{h} + \nabla \tilde{G} = 0 \quad \text{and} \quad \tilde{G}(x) = \tilde{Z}(x) \tilde{F}(x), \tag{2.56}
\]

where

\[
\tilde{u} = e^{i\tau} u, \quad \tilde{F} = e^{i\tau} F, \quad \tilde{G} = e^{i(\tau + \theta)} G, \quad \tilde{h} = e^{i(\tau + \theta)} h \quad \text{and} \quad \tilde{Z}(x) = e^{i\theta} Z(x), \tag{2.57}
\]

in which \(\tau\) and \(\theta\) are real parameters. Therefore, the fields with tildes on them satisfy the same equations as those without tildes, and the variational principles directly apply provided \(\theta\) is chosen so that

\[
\tilde{Z}''(x) \geq \alpha I, \quad \forall x \in \Gamma, \tag{2.58}
\]

for some \(\alpha > 0\). With the subsequent replacement (2.57), we obtain a new set of variational principles parametrized by \(\tau\) and \(\theta\). Observe that even if (2.15) is not satisfied for any \(\alpha > 0\), we still may be able to find a range of values of \(\theta\) for which (2.58) holds, and thus for which we can obtain variational principles. With the particular choice \(\theta = 0\) and \(e^{i\tau} = -i\), the variational principle (2.41) gets mapped to the dual variational principle (2.51).

All the previous minimization principles need the fundamental coercivity assumption (2.15). If \(Z''(x) = 0\) in some part \(\Psi\) of the domain \(\Gamma\), they cannot be applied directly. However, we can still obtain variational principles by taking appropriate limits. If, in some region \(\Psi \subset \Gamma\), \(Z''(x)\) is very small in the sense that

\[
\delta I \geq Z''(x), \quad \forall x \in \Psi, \tag{2.59}
\]

for some very small \(\delta > 0\), then within this region,

\[
\left( \frac{\nabla u'}{-G'} \right) \cdot \mathcal{L} \left( \frac{\nabla u'}{-G'} \right) \approx [Z' \cap u' - G'] \cdot (Z'')^{-1} [Z' \cap u' - G']. \tag{2.60}
\]

If \(\Psi = \Gamma\), then the variational principle reduces to a variational principle similar to (1.12), which is not particularly interesting. If \(\Psi \not\subset \Gamma\), the variational principle will only be useful if one takes \(G' \approx Z' \cap u'\) in \(\Psi\). The variational principles still

Proc. R. Soc. A (2009)
hold in the limit as \( \delta \to 0 \) (but should be used carefully because the lack of coercivity in the region \( \Psi \) does not enable one to state existence and uniqueness of a solution in this case) provided one first takes
\[
\mathcal{G}' = Z' \cap u', \quad \forall x \in \Psi,
\] (2.61)
corresponding to taking an exact solution of (2.11) in this region (but with \( (u', \mathcal{G}') \) not necessarily coinciding with the exact solution \( (u', \mathcal{G}') \) associated with the given boundary conditions). Then, as \( \delta \to 0 \), the expression for \( Y \) reduces to
\[
Y(u', \mathcal{G}') = \int_{\Gamma \setminus \Psi} \left( \frac{\nabla u'}{-\mathcal{G}'} \right) \cdot \mathcal{L} \left( \frac{\nabla u'}{-\mathcal{G}'} \right) - \int_{\Gamma} 2 h'' \cdot u'.
\] (2.62)
Outside \( \Psi \), we still require that \( Z''(x) \geq \alpha I \), for some \( \alpha > 0 \).

We now show how these variational principles may give useful information about the Dirichlet to Neumann map on the boundary \( \partial \Gamma \). Let us assume \( h = 0 \) and introduce the \( m \)-component complex vector fields
\[
q_0 = G(x) \cdot n \quad \text{and} \quad u_0(x) = u(x), \quad \text{for} \ x \in \partial \Gamma.
\] (2.63)
The relation between \( q_0 \) and \( u_0 \) must be linear and we can write \( q_0 = Nu_0 \) that defines the Dirichlet to Neumann map \( N \). If \( u_0^{(1)} \) and \( u_0^{(2)} \) denote two boundary conditions for \( u \) and \( G^{(j)}, \mathcal{F}^{(j)} \) and \( q_0^{(j)}, j = 1, 2 \), denote the associated fields, then it follows from the key property and the symmetry of \( Z(x) \) that
\[
\left\langle Nu_0^{(1)}, u_0^{(2)} \right\rangle = \int_{\partial \Gamma} q_0^{(1)} \cdot u_0^{(2)} = \int_{\Gamma} G^{(1)} \cdot \mathcal{F}^{(2)} = \int_{\Gamma} \mathcal{F}^{(1)} \cdot Z \mathcal{F}^{(2)}
\]
\[
= \int_{\Gamma} G^{(2)} \cdot \mathcal{F}^{(1)} = \int_{\partial \Gamma} q_0^{(2)} \cdot u_0^{(1)} = \left\langle u_0^{(1)}, Nu_0^{(2)} \right\rangle,
\] (2.64)
which implies that the map \( N \) is symmetric. It has a positive semi-definite imaginary part since
\[
\left\langle u_0', N''u_0' \right\rangle + \left\langle u_0'', Nu''u_0'' \right\rangle = \int_{\partial \Gamma} q''_0 \cdot u_0' - q''_0 \cdot u_0' = \int_{\Gamma} G'' \cdot \mathcal{F}' - G' \cdot \mathcal{F}''
\]
\[
= \int_{\Gamma} \mathcal{F}' \cdot Z'' \mathcal{F}' + \mathcal{F}'' \cdot Z'' \mathcal{F}'' \geq 0.
\] (2.65)
Assuming \( N'' \) is invertible, we can write the relation \( q_0 = Nu_0 \) as
\[
\begin{pmatrix} q''_0 \\ u''_0 \end{pmatrix} = \mathcal{N} \begin{pmatrix} u'_0 \\ -q'_0 \end{pmatrix},
\] (2.66)
where
\[
\mathcal{N} = \begin{pmatrix} N'' + N'(N'')^{-1}N' & N'(N'')^{-1} \\ (N'')^{-1}N' & (N'')^{-1} \end{pmatrix}.
\] (2.67)
Thus, the map \( \mathcal{N} \) giving the boundary fields \( q''_0 \) and \( u''_0 \) in terms of \( q'_0 \) and \( u'_0 \) must have this special form.
Now, using the key property, \( Y(u', G') \) can be expressed (when \( h = 0 \)) in terms of the boundary fields

\[
Y(u', G') = \int_G G'' \cdot F' - G' \cdot F'' = \int_{\partial G} q''_0 \cdot u'_0 - q'_0 \cdot u''_0 = \left\langle \begin{pmatrix} u'_0 \\ -q'_0 \end{pmatrix}, N \begin{pmatrix} u'_0 \\ -q'_0 \end{pmatrix} \right\rangle,
\]

and thus bounds on \( Y(u', G') \), obtained from the variational principle, give upper bounds on the quadratic form associated with \( N \). Bounds on \( Y(u'', G'') \), obtained from the dual variational principle, give lower bounds on the quadratic form associated with \( N \).

### 3. Application to acoustics and elastodynamics

Let us begin by studying the acoustic equation in \( d \) dimensions,

\[
-\nabla \cdot \rho^{-1} \nabla P = (1/\kappa) \omega^2 P,
\]

where \( P(x) \) is the complex pressure; \( \kappa(x) \) is the bulk modulus; \( \rho(x) \) is the density tensor; and \( \omega \) is the (fixed) frequency of oscillation. (As is well known, this equation with \( d = 2 \) is also applicable in cylindrical bodies, for antiplane elastodynamics, where the displacement field is directed parallel to the cylinder axis, and for electromagnetism when the fields are transverse electric (TE), where the electric field is directed parallel to the cylinder axis, or transverse magnetic (TM), where the magnetic field is directed parallel to the cylinder axis.) Normally, one expects \( \kappa \) to have a negative imaginary part (due to bulk viscosity) and \( \rho \) to be \( \rho \mathbb{I} \), where \( \rho(x) \) is the positive mass density. However, viewed as the effective density tensor of a composite material at a given frequency \( \omega \), \( \rho \) can be anisotropic (Schoenberg & Sen 1983) and, at least in the context of elastodynamics, can be complex valued with a positive semi-definite imaginary part (Milton & Willis 2007) and with a real part that is not even necessarily positive (as established by Bouchitté & Felbacq 2004, Movchan & Guenneau 2004 and Felbacq & Bouchitté 2005 in the context of antiplane elastodynamics, where \( 1/\kappa \) plays the role of density, and by Ávila et al. (2005) and Liu et al. (2005) in the context of three-dimensional elastodynamics). Such unusual effective densities are due to the fact that different parts of the microstructure can experience different accelerations, and can be out of phase with the overall applied force if this relative motion dissipates energy.

Given the complex pressure field \( P(x) \), the associated complex velocity field of the fluid is

\[
v = -i(\omega \rho)^{-1} \nabla P.
\]

(The associated physical pressure and physical velocity are, respectively, \( (e^{-i\omega t} P(x))' \) and \( (e^{-i\omega t} v(x))' \), where the prime denotes the real part and \( t \) denotes the time.) Upon multiplying (3.1) by \( e^{i\theta} \) and comparing it with (2.12), we can make the identifications

\[
u = P, \quad L = -e^{i\theta} \rho^{-1}, \quad K = 0, \quad M = e^{i\theta} \omega^2/\kappa, \quad h = 0, \quad G = -i e^{i\theta} \omega v.
\]

\[\text{Proc. R. Soc. A (2009)}\]
Thus, the potential $u$ is a scalar, corresponding to the case $m=1$. We choose $\theta$ such that the imaginary part of

$$Z = \begin{pmatrix} -e^{i\theta} \rho^{-1} & 0 \\ 0 & e^{i\theta} \omega^2 / \kappa \end{pmatrix}$$

satisfies (2.15). When $\rho$ is real and positive definite, as is typically the case, then we need to choose $\theta$ in the range $0 > \theta > -\pi$. If, in addition, $\kappa$ is almost real and positive, with a small negative imaginary part, then we should choose $\theta$ appropriately small and negative.

Explicit expressions for the various variational principles in the acoustic case can of course be obtained by making the substitutions (3.3) in the relevant equations in §2. For example, if $\rho$ has a positive definite imaginary part (which is more applicable in the mathematically analogous TE or TM electromagnetic problems) and $\kappa(x)$ has a negative imaginary part, then, with $\theta=0$, we have the variational principle

$$Y(P', u') = \inf_{P'} \inf_{u'} Y(P', u'),$$

where

$$Y(P', u') = \int \left( \begin{array}{c} \nabla P' \\ -\omega u' \end{array} \right) \cdot \mathcal{R} \left( \begin{array}{c} \nabla P' \\ -\omega u' \end{array} \right) + \left( \begin{array}{c} \omega P' \\ -\nabla \cdot u' \end{array} \right) \cdot \mathcal{K} \left( \begin{array}{c} \omega P' \\ -\nabla \cdot u' \end{array} \right),$$

in which

$$\mathcal{R} = \begin{pmatrix} r'' + r'(r'')^{-1}r' & r'(r'')^{-1} \\ (r'')^{-1}r' & (r'')^{-1} \end{pmatrix} \quad \text{and} \quad \mathcal{K} = \begin{pmatrix} k'' + (k')^2 / k'' & k' / k'' \\ k' / k'' & 1 / k'' \end{pmatrix},$$

where $r = -\rho^{-1}$ and $k = 1 / \kappa$. In this variational principle, the boundary values of $P'$ and $u' \cdot n$ are fixed, and the infimums are attained at fields $P'$ and $u''$ associated with solutions of the acoustic equation satisfying the prescribed boundary conditions. That the Euler–Lagrange equations of (3.5) coincide with (3.1) and (3.2) follows from the analysis of (2.49) and (2.50), and can also be shown directly.

When $r''$ is very small within all of $\Gamma$, then we have

$$\left( \begin{array}{c} \nabla P' \\ -\omega u' \end{array} \right) \cdot \mathcal{R} \left( \begin{array}{c} \nabla P' \\ -\omega u' \end{array} \right) \approx \left[ r' \nabla P' - \omega u' \right] \cdot \left( r'' \right)^{-1} \left[ r' \nabla P' - \omega u' \right],$$

and the variational principles will only be useful if we take trial fields with $\omega u'' \approx r' \nabla P'$. The variational principles still hold in the limit as $r'' \to 0$, provided we first choose $u' = r' \nabla P' / \omega$. Then, when $\rho$ is real, we have the variational principle

$$Y(P') = \inf_{P'} Y(P'),$$

where

$$Y(P') = \int \left( \begin{array}{c} \omega P' \\ -\rho^{-1} \nabla P' / \omega \end{array} \right) \cdot \mathcal{K} \left( \begin{array}{c} \omega P' \\ -\rho^{-1} \nabla P' / \omega \end{array} \right).$$
In this variational principle, the boundary values of \( P' \) and \( \mathbf{n} \cdot \rho^{-1}\nabla P \) are fixed, where the latter corresponds to fixing the boundary value of \( \psi'' \cdot \mathbf{n} \). If \( \rho \) is discontinuous across an interior interface, then the trial field \( P' \) should be chosen so that the normal component of \( \rho^{-1}\nabla P' \) is continuous across this interface.

The infimum is attained at a field \( P' \) associated with solutions of the acoustic equation satisfying the prescribed boundary conditions. To show this, let \( P' \) be a minimizer and set \( \rho^{-1}\nabla P' = \mathbf{n} \cdot \rho^{-1}\nabla S = 0 \) on \( \partial I \). A necessary condition for the minimum to be at \( S = 0 \) is that the first derivative of the functional \( Y \) vanishes, i.e.

\[
\int_I \left( U \left( \omega P'' \right) \cdot \left( \omega S \right) \right) = 0, \tag{3.11}
\]

where we have introduced

\[
\begin{pmatrix}
U \\
\omega P''
\end{pmatrix} = \mathcal{K} \begin{pmatrix}
\omega S \\
\nabla \cdot (\rho^{-1}\nabla P'/\omega)
\end{pmatrix}. \tag{3.12}
\]

Integrating constraint (3.11) by parts gives

\[
\int_I S(\omega U + \nabla \cdot (\rho^{-1}\nabla P')) = 0, \tag{3.13}
\]

and this will be satisfied for any \( S \), with \( S = 0 \) and \( \mathbf{n} \cdot \rho^{-1}\nabla S = 0 \) on \( \partial I \), if and only if

\[
U = -\nabla \cdot (\rho^{-1}\nabla P''/\omega). \tag{3.14}
\]

A straightforward calculation (similar to the equivalence of (2.7) and (2.35)) shows that (3.12) is equivalent to (3.1) with this value of \( U \). Thus, the acoustic equation is the Euler–Lagrange equation of (3.9), when \( P'' \) is defined by (3.12).

This variational principle remains valid if we replace \( P \) by \( e^{\mathbf{i} \tau} P \), where \( \tau \) can be any constant.

In a similar way, if \( \rho \) has a positive definite imaginary part but \( \kappa \) is real (which is more applicable to the analogous TM electromagnetic problem, where the magnetic permeability plays the role of \( 1/\kappa \)), then we have the variational principle

\[
Y(\psi'') = \inf_{\psi''} Y(\psi''), \tag{3.15}
\]

where

\[
Y(\psi'') = \int_I \left( \nabla(\kappa \nabla \cdot \psi'')/\omega \right) \cdot \mathcal{R} \left( \nabla(\kappa \nabla \cdot \psi'')/\omega \right), \tag{3.16}
\]

in which the boundary values of \( \psi'' \cdot \mathbf{n} \) and \( \kappa \nabla \cdot \psi'' \) are fixed, where the latter corresponds to fixing the boundary value of \( P' \).

One has to slightly modify the analysis to apply it to the elastodynamic equations

\[
-\nabla \cdot (C \nabla \mathbf{u}) = b + \omega^2 \rho \mathbf{u}, \tag{3.17}
\]

in which \( \mathbf{u}(\mathbf{x}) \) is the complex displacement field; \( b(\mathbf{x}) \) is the complex body-force density; \( \rho(\mathbf{x}) \) is the complex density tensor; and \( C(\mathbf{x}) \) is the complex elasticity tensor. Given the complex displacement field \( \mathbf{u}(\mathbf{x}) \), the complex stress \( \mathbf{t}(\mathbf{x}) \) and
complex momentum density \( p(x) \) are given by

\[
\tau = C \nabla u \quad \text{and} \quad p = -i \omega \rho u. \tag{3.18}
\]

(The associated physical displacement, body-force density, stress and momentum density are, respectively, \( (e^{-i\omega t} u)' \), \( (e^{-i\omega t} b)' \), \( (e^{-i\omega t} \tau)' \) and \( (e^{-i\omega t} p)' \).) Upon multiplying (3.17) by \( e^{i\theta} \) and comparing it with (2.12), we can make the identifications

\[
L = -e^{i\theta} C, \quad K = 0, \quad M = e^{i\theta} \omega^2 \rho \quad \text{and} \quad h = e^{i\theta} b, \tag{3.19}
\]

and (3.18) implies

\[
G = -e^{i\theta} \tau \quad \text{and} \quad g = i e^{i\theta} \omega p. \tag{3.20}
\]

The problem is that \( C \) acting upon any antisymmetric matrix is zero, and so \( L \) and hence \( Z \) are singular, which is a problem for computing the inverses in equation (2.36) for \( L \). This problem is rectified in the standard way, by replacing \( \square \) everywhere with \( \square \) defined by

\[
\square u \equiv \left( \frac{[\nabla u + (\nabla u)^T]/2}{u} \right).
\tag{3.21}
\]

Then, the fields \( F \) and \( G \) have the form (2.1), with \( F \) and \( G \) being symmetric \( d \times d \) matrices; thus, \( F \) and \( G \) can be represented by \( d(d+3)/2 \)-component vectors. We choose \( \theta \) such that the imaginary part of \( Z \) satisfies (2.15) on the space of these fields, and we compute the inverse of \( Z'' \) on the space of these fields. When \( \rho \) is real and positive definite, then we need to choose \( \theta \) in the range \( \pi > \theta > 0 \). If, in addition, \( C \) is almost real with a positive definite real part, and with a small negative definite imaginary part, then we should choose \( \theta \) appropriately small and positive. With \( \square \) replaced by \( \square \), the analysis of equations (2.23)–(2.29) and (2.44)–(2.50) still holds, and we recover the elastodynamic equations from the Euler–Lagrange equations associated with the variational principles.

In the special case \( \theta = 0 \), the functionals \( Q, R \) and \( Y \) reduce to

\[
Q(u', u'') = \int_I \left( -\epsilon \cdot C \epsilon + \omega^2 u \cdot \rho u + 2 b \cdot u \right)'',
\]

\[
R(\tau', \tau'') = \int_I \left[ -\tau \cdot C^{-1} \tau + (\nabla \cdot \tau + b) \cdot (\omega^2 \rho)^{-1} (\nabla \cdot \tau + b) \right]''
\]

and

\[
Y(u', \tau') = \int_I - \left( \begin{array}{c} \epsilon' \\ -\tau' \end{array} \right) \cdot C \left( \begin{array}{c} \epsilon' \\ -\tau' \end{array} \right) \right) \]

\[
+ \left( \begin{array}{c} \omega u' \\ (b' + \nabla \cdot \tau')/\omega \end{array} \right) \mathcal{P} \left( \begin{array}{c} \omega u' \\ (b' + \nabla \cdot \tau')/\omega \end{array} \right) + 2 b'' \cdot u',
\tag{3.23}
\]

\( \text{Proc. R. Soc. A} \ (2009) \)
where

\[
\epsilon = \frac{[\nabla u + (\nabla u)^T]}{2},
\]

\[
C = \begin{pmatrix}
C'' + C'(C'')^{-1} C' & C'(C'')^{-1} \\
(C'')^{-1} C' & (C'')^{-1}
\end{pmatrix},
\]

\[
P = \begin{pmatrix}
\rho'' + \rho'(\rho'')^{-1} \rho' & \rho'(\rho'')^{-1} \\
(\rho'')^{-1} \rho' & (\rho'')^{-1}
\end{pmatrix}
\]

are, respectively, negative definite and positive definite matrices. It follows directly from the variational principles of Gurtin (1964) and Willis (1981), applied to the time-harmonic case considered here, that the stationary points of the functionals \(Q\) and \(R\) correspond to solutions of the elastodynamics equations, in agreement with the results of §2.

When \(r_0\) is very small, we have

\[
\frac{\omega u'}{(b' + \nabla \cdot \tau')/\omega) P \left( \frac{\omega u'}{(b' + \nabla \cdot \tau')/\omega} \right) \approx \left[ \frac{\omega^2 \rho' u' + b' + \nabla \cdot \tau'}{(\omega^2 \rho'')^{-1}} \right] X \left[ \frac{\omega^2 \rho' u' + b' + \nabla \cdot \tau'}{(\omega^2 \rho'')^{-1}} \right],
\]

and the variational principles will only be useful if we take trial fields with \(\omega^2 \rho' u' + b' + \nabla \cdot \tau'\) approximately zero. The variational principle involving the functional \(Y\) still holds in the limit as \(\rho'' \to 0\), provided we first choose

\[
u' = -(\rho')^{-1}(b' + \nabla \cdot \tau')/\omega^2.
\]

Then, when \(\rho\) is real, we have the variational principle

\[
Y(\tau') = \inf_{\tau'} Y(\tau'),
\]

where

\[
Y(\tau') = \int_{\Gamma} \left( \frac{\epsilon'}{-\tau'} - \epsilon' \right) \cdot C \left( \frac{\epsilon'}{-\tau'} \right) + 2b'' \cdot u',
\]

in which \(u'\) and \(\epsilon'\) are given by (3.27) and (3.24). In this variational principle, the trial field \(\tau'\) can be any real symmetric matrix-valued field, with prescribed values of \(\tau' \cdot n\) and \(\rho^{-1} \nabla \cdot \tau'\) at the boundary \(\partial \Gamma\), where prescribing \(\rho^{-1} \nabla \cdot \tau'\) at \(\partial \Gamma\) corresponds to prescribing \(\u'\) according to (3.27). The infimum is attained at a field \(\tau'\) associated with solutions of the elastodynamic equation satisfying the prescribed boundary conditions. If \(\rho\) is discontinuous across an interior interface, then the trial field \(\tau'\) needs to be chosen so that \(u'\) given by (3.27) is continuous across this interface.
4. The saddle-point variational principles for electromagnetism

A comparison of the continuum elastodynamic equation (3.17) with Maxwell’s equations,

\[ \nabla \times E = i\omega B, \quad \nabla \times H = -i\omega D + j, \quad D = \varepsilon E \quad \text{and} \quad B = \mu H, \quad (4.1) \]

which can be rewritten (Milton et al. 2006) in the form

\[ -\nabla \cdot (C \nabla E) = i\omega J + \omega^2 \varepsilon E, \quad (4.2) \]

where now \( C \) has elements

\[ C_{ijk} = e_{ijm} e_{kln} \{ \mu^{-1} \}_{mn}, \quad (4.3) \]

in which \( E \) is the electric field; \( J \) is the free current (not including conduction currents \( \sigma E \), which are instead incorporated in the term \( \varepsilon E \)); \( \varepsilon \) is the electric permittivity tensor; \( \mu \) is the magnetic permeability tensor; and \( e_{ijm} \) is the Levi-Civita (alternating) tensor \((e_{ijm} = 1 \ (-1) \text{ if } (i, j, m) \text{ is an even (odd) permutation of } (1, 2, 3) \text{ and is zero otherwise})\), suggests that the preceding analysis should also extend to three-dimensional electromagnetism. Upon multiplying (4.2) by \( e^{i\theta} \) and comparing it with (2.12), we can make the identifications

\[ \begin{align*}
  u &= E, \\
  L &= -e^{i\theta} C, \\
  K &= 0, \\
  M &= e^{i\theta} \omega^2 \varepsilon \\
  h &= i\omega e^{i\theta} J. 
\end{align*} \quad (4.4) \]

The problem is now that \( C \) acting upon any symmetric matrix is zero, and so \( L \) and hence \( Z \) are singular, which is a problem for computing the inverses in equation (2.36) for \( L \). This problem is rectified by replacing \( \nabla \) everywhere with \( \hat{\nabla} \) defined by

\[ \hat{\nabla} E = \left( \frac{[\nabla E - (\nabla E)^\top]/2}{E} \right). \quad (4.5) \]

Then, the fields \( F \) and \( G \) have the form (2.1), with \( F \) and \( G \) being antisymmetric \( 3 \times 3 \) matrices; thus, \( F \) and \( G \) can be represented by six component vectors. We choose \( \theta \) such that the imaginary part of

\[ Z = \begin{pmatrix}
  -e^{i\theta} C & 0 \\
  0 & e^{i\theta} \omega^2 \varepsilon
\end{pmatrix} \quad (4.6) \]

satisfies (2.15) on the space of these fields, and we compute the inverse of \( Z'' \) on the space of these fields.

Let us assume, for simplicity, that

\[ \varepsilon''(x) \geq \alpha_1 I \quad \text{and} \quad \mu''(x) \geq \alpha_2 I, \quad \forall x \in \Gamma, \quad (4.7) \]

for some \( \alpha_1 > 0 \) and \( \alpha_2 > 0 \). (By multiplying (4.2) by \( e^{i\theta} \) and redefining \( \varepsilon, \mu \) and \( J \) if necessary.) Then, we can take \( \theta = 0 \) and (2.15) holds for some \( \alpha > 0 \). From (4.1) and (4.5), we note that

\[ F = \hat{\nabla} E = \begin{pmatrix} F \\ E \end{pmatrix}, \quad \text{with} \quad F = \frac{i\omega}{2} \begin{pmatrix} 0 & -B_3 & B_2 \\
  B_3 & 0 & -B_1 \\
 -B_2 & B_1 & 0 \end{pmatrix}, \quad (4.8) \]

Proc. R. Soc. A (2009)
and the constitutive law (2.9) implies

\[
\mathcal{G} = \begin{pmatrix} G \\ \omega^2 D \end{pmatrix}, \quad \text{with } G = -i\omega \begin{pmatrix} 0 & -H_3 & H_2 \\ H_3 & 0 & -H_1 \\ -H_2 & H_1 & 0 \end{pmatrix}.
\]

Thus, the key property (2.4) reduces to

\[
\int \omega^2 (B \cdot H + E \cdot D) + i\omega j \cdot E = \int_{\partial \Gamma} -i\omega (H \times n) \cdot E = \int_{\partial \Gamma} -i\omega (E \times H) \cdot n,
\]

and holds for all fields \( E, D, B \) and \( H \) satisfying the differential constraints

\[
\nabla \times E = i\omega B \quad \text{and} \quad \nabla \times H = -i\omega D + j.
\]

Now the first relation (2.20) will hold if

\[
\int_{\partial \Gamma} (G'' \cdot n) \cdot s = \int_{\partial \Gamma} -i\omega (H'' \times n) \cdot s = 0,
\]

and for this it suffices that the tangential component of \( s \) vanishes at the surface \( \partial \Gamma \). Similarly, the second relation in (2.20) will hold if this condition is satisfied. Therefore, the infimum and supremum in (2.17) can be extended to all \( u' \) and \( u'' \), with the same tangential components as, respectively, \( E' \) and \( E'' \). Similarly, in the minimization variational principle (2.41) or (2.51), we need only to require that \( u' \) or \( u'' \) has the same tangential components as \( E' \) or \( E'' \), respectively.

With \( \cap \) replaced by \( \cap \), the analysis of equations (2.23)–(2.29) and (2.44)–(2.50) still holds, and we recover Maxwell’s equations from the Euler–Lagrange equations associated with the variational principles. If we introduce the tangential components \( E_\parallel(x) = E - n (n \cdot E) \) and \( H_\parallel(x) = H - n (n \cdot H) \), then

\[
q_0(x) = G \cdot n = -i\omega (H_\parallel \times n).
\]

When \( j=0 \), we can write \( q_0 = NE_\parallel \), thus defining the map \( N \) that governs the electrodynamic response of the body at the frequency \( \omega \). The analysis of (2.63)–(2.68) shows that the map \( N \) is symmetric with a positive semi-definite imaginary part, and can be bounded (when \( j=0 \)) using the minimization variational principles.

Equivalently, we can find the variational principles directly from Maxwell’s equations (4.1). The key property (basically Poynting’s theorem) (4.10) is obtained by integrating over \( \Gamma \) the identity

\[
\nabla \cdot (E \times H) = H \cdot (\nabla \times E) - E \cdot (\nabla \times H) = i\omega (H \cdot B + E \cdot D) - j \cdot E.
\]

Motivated by the form of (4.8) and (4.9), let us redefine the fields

\[
\mathcal{F} = \begin{pmatrix} F \\ f \end{pmatrix} = \begin{pmatrix} i\omega B \\ E \end{pmatrix} \quad \text{and} \quad \mathcal{G} = \begin{pmatrix} G \\ g \end{pmatrix} = \begin{pmatrix} -i\omega H \\ \omega^2 D \end{pmatrix}.
\]
which are subject to the differential constraints that

\[ F = \Omega E \quad \text{and} \quad i\omega j + \overline{\sigma} G = 0, \]  

(4.17)

where the operators \( \Omega \) and \( \overline{\sigma} \) are defined by

\[ \Omega E = \left( \nabla \times \frac{E}{E} \right) \quad \text{and} \quad \overline{\sigma} G = \nabla \times G + g, \]  

(4.18)

so that the differential constraints (4.11) and (4.17) are equivalent. The key property (4.10) now takes the form

\[ \int_{\Gamma} G \cdot F + i\omega j \cdot E = \int_{\partial \Gamma} (E \times G) \cdot n, \]  

(4.19)

where \( G \cdot F = F' \cdot G + f \cdot g \). The constitutive law \( G = ZF \) holds with \( Z \) redefined as

\[ Z = \begin{pmatrix} -\mu^{-1} & 0 \\ 0 & \omega^2 \varepsilon \end{pmatrix}, \]  

(4.20)

and this can be re-expressed in the form (2.14), with

\[ F' = \Omega E' = \left( \begin{array}{c} -\omega B'' \\ E' \end{array} \right), \quad F'' = \Omega E'' = \left( \begin{array}{c} \omega B'' \\ E'' \end{array} \right), \]  

\[ G' = \left( \begin{array}{c} \omega H'' \\ \omega^2 D' \end{array} \right) \quad \text{and} \quad G'' = \left( \begin{array}{c} -\omega H'' \\ \omega^2 D' \end{array} \right), \]  

(4.21)

where these fields satisfy the differential constraints

\[ F' = \Omega E', \quad -\omega j'' + \overline{\sigma} G' = 0, \quad F'' = \Omega E'' \quad \text{and} \quad \omega j' + \overline{\sigma} G'' = 0, \]  

(4.22)

implying, as a corollary of (4.19), that

\[ \int_{\Gamma} G' \cdot \Omega s - \omega j'' \cdot s = 0 \quad \text{and} \quad \int_{\Gamma} G'' \cdot \Omega s + \omega j' \cdot s = 0, \]  

(4.23)

for all real-valued vector fields \( s \), such that the tangential component of \( s \) vanishes at the surface \( \partial \Gamma \). Introducing the functional

\[ Q(E', E'') = \int_{\Gamma} \left( \frac{\Omega E'}{\Omega E''} \right) \cdot \left( \begin{array}{cc} Z'' & Z' \\ Z' & Z'' \end{array} \right) \left( \frac{\Omega E'}{\Omega E''} \right) + 2\omega (j' \cdot E' - j'' \cdot E'') \]  

\[ = \int_{\Gamma} [(\Omega E) \cdot Z \Omega E + 2i\omega j \cdot E]'' \]  

(4.24)

we note from (4.23) that

\[ Q(E' + s, E'') = Q(E', E'') + \int_{\Gamma} (\Omega s) \cdot Z'' \Omega s, \]  

\[ Q(E', E'' + s) = Q(E', E'') - \int_{\Gamma} (\Omega s) \cdot Z'' \Omega s. \]  

(4.25)
Consequently, the pair of fields \((\mathbf{E}', \mathbf{E}'') = (\mathbf{E}', \mathbf{E}'')\) is at the saddle point of (2.17) and (2.18), where the infimum and supremum are over all vector fields \(\mathbf{E}'\) and \(\mathbf{E}''\), with the same tangential components at the boundary \(\partial \Gamma\) as, respectively, \(\mathbf{E}'\) and \(\mathbf{E}''\). Owing to the diagonal structure (4.20) of the matrix \(\mathbf{Z}(\mathbf{x})\), the formula (4.24) for \(Q(\mathbf{E}', \mathbf{E}'')\) reduces to

\[
Q(\mathbf{E}', \mathbf{E}'') = \int_{\Gamma} \left[ (\nabla \times \mathbf{E}) \cdot \mathbf{m} (\nabla \times \mathbf{E}) + \omega^2 \mathbf{E} \cdot \mathbf{e} \mathbf{E} + 2i\omega \mathbf{j} \cdot \mathbf{E} \right],
\]

where \(\mathbf{m} = -\mathbf{\mu}^{-1}\), as with \(\mathbf{e}\), has a positive definite imaginary part. Hence, we have the saddle-point variational principles

\[
Q(\mathbf{E}', \mathbf{E}'') = \inf_{\mathbf{E}'} \sup_{\mathbf{E}''} Q(\mathbf{E}', \mathbf{E}'') = \sup_{\mathbf{E}'} \inf_{\mathbf{E}''} Q(\mathbf{E}', \mathbf{E}''),
\]

where

\[
Q(\mathbf{E}', \mathbf{E}'') = \int_{\Gamma} \mathbf{G}'' \cdot \mathbf{F}' + \mathbf{G}' \cdot \mathbf{F}'' + 2\omega (\mathbf{j} \cdot \mathbf{E}' - \mathbf{j}'' \cdot \mathbf{E}'')
\]

\[= \int_{\Gamma} \omega^2 (\mathbf{H}' \cdot \mathbf{B}'' + \mathbf{H}'' \cdot \mathbf{B}' + \mathbf{E}' \cdot \mathbf{D}'' + \mathbf{E}'' \cdot \mathbf{D}') + 2\omega (\mathbf{j} \cdot \mathbf{E}' - \mathbf{j}'' \cdot \mathbf{E}').
\]

When \(\mathbf{j} = 0\), we can use (4.15) to express \(Q(\mathbf{E}', \mathbf{E}'')\) just in terms of the tangential components of the fields at the boundary,

\[
Q(\mathbf{E}', \mathbf{E}'') = \omega \int_{\partial \Gamma} (\mathbf{E}'' \times \mathbf{H}'') \cdot \mathbf{n} - (\mathbf{E}' \times \mathbf{H}') \cdot \mathbf{n}, \quad \text{when } \mathbf{j} = 0.
\]

The result that the stationary point of the variational principle (4.27) corresponds to a solution of Maxwell’s equations follows from a variational principle derived by Willis (1984). The Willis variational principle is expressed in terms of the electromagnetic potentials \(V(\mathbf{x})\) and \(\mathbf{A}(\mathbf{x})\), which are associated with \(\mathbf{E}\) and \(\mathbf{B}\) through the relations

\[
\mathbf{E} = -\nabla V + i\omega \mathbf{A} \quad \text{and} \quad \mathbf{B} = \nabla \times \mathbf{A},
\]

and, in the time-harmonic case, involves (within a proportionality constant) the functional

\[
\hat{Q}(V, A) = \omega^2 \int_{\partial \Gamma} \left\{ [-((\nabla \times \mathbf{A}) \cdot \mathbf{m} (\nabla \times \mathbf{A}) + (\nabla V - i\omega \mathbf{A}) \cdot \mathbf{e} (\nabla V - i\omega \mathbf{A})
\]

\[+ 2i\nabla \cdot \mathbf{j}/\omega - 2\mathbf{A} \cdot \mathbf{j}] e^{-i\omega t}\right\},
\]

where, as usual, the prime denotes the real part. His variational principle states that, for any choice of \(t\), the functional \(\hat{Q}(V, A)\) applied to trial fields \(V\) and \(\mathbf{A}\), having any given fixed values at the boundary \(\partial \Gamma\), has a stationary point at the potentials \(V\) and \(\mathbf{A}\) solving Maxwell’s equations, and having the prescribed boundary values. Setting \(\mathbf{E} = -\nabla V + i\omega \mathbf{A}\), one notes by integrating by parts that \(\hat{Q}(V, A)\) is just a functional of \(\mathbf{E}\) alone, and when \(t\) is chosen, with \(e^{-i\omega t} = -i\), one has the identity

\[
\hat{Q}(V, A) = Q(\mathbf{E}', \mathbf{E}'').
\]

Furthermore, the values of \(V\) and \(\mathbf{A}\) on \(\partial \Gamma\) determine the tangential (but not normal) components of \(\mathbf{E}'\) and \(\mathbf{E}''\) on \(\partial \Gamma\).

Proc. R. Soc. A (2009)
There is also a dual variational principle. By taking real and imaginary parts of the constitutive relation \( \mathcal{F} = \mathbf{K} \mathcal{G} \), where \( \mathbf{K} = \mathbf{Z}^{-1} \) has a negative definite imaginary part, we obtain (2.23). Let us introduce the functional

\[
R(\mathbf{H}', \mathbf{H}'') = \int \left( \mathcal{G}' \left( \mathbf{K}'' \mathbf{K}' - \mathcal{G}'' \right) \mathcal{G}' \right) = \int (\mathcal{G} \cdot \mathbf{K} \mathcal{G})'',
\]

where \( \mathcal{G}' \) and \( \mathcal{G}'' \) are given by

\[
\mathcal{G}' = \omega \left( \frac{\mathbf{H}''}{j'' - \nabla \times \mathbf{H}''} \right) \quad \text{and} \quad \mathcal{G}'' = \omega \left( \frac{-\mathbf{H}'}{-j' + \nabla \times \mathbf{H}'} \right)
\]

(to ensure that \(-\omega j'' + \mathcal{G}' = 0\) and \(\omega j' + \mathcal{G}'' = 0\) inside \(\Gamma\)), and \(\mathbf{H}''\) and \(\mathbf{H}'\) have the same tangential components as, respectively, \(\mathbf{H}''\) and \(\mathbf{H}'\). We then have the saddle-point variational principles

\[
R(\mathbf{H}', \mathbf{H}'') = \sup_{\mathbf{H}''} \inf_{\mathbf{H}'} R(\mathbf{H}', \mathbf{H}''),
\]

where the supremum and infimum are over fields with the required tangential components at the boundary. Owing to the diagonal structure (4.20) of \(\mathbf{Z}(\mathbf{x})\), the formula for \(R(\mathbf{H}', \mathbf{H}'')\) reduces to

\[
R(\mathbf{H}', \mathbf{H}'') = \int \left[ (j - \nabla \times \mathbf{H}) \cdot (j - \nabla \times \mathbf{H}) + \omega^2 \mathbf{H} \mathbf{\mu} \mathbf{H} \right],
\]

where \(e = -\varepsilon^{-1}\), as with \(\mathbf{\mu}\), has a positive definite imaginary part. From this expression, one can see that \(R(\mathbf{H}', \mathbf{H}'')\) is a special case of the functional associated with the Willis (1984) dual variational principle. We also have that

\[
R(\mathbf{H}', \mathbf{H}'') = \int \mathcal{G}'' \cdot \mathcal{F}' + \mathcal{G}' \cdot \mathcal{F}'' = \int \omega^2 (\mathbf{H}' \cdot \mathbf{B}'' + \mathbf{H}'' \cdot \mathbf{B}' + \mathbf{E}' \cdot \mathbf{D}'' + \mathbf{E}'' \cdot \mathbf{D}'),
\]

which, when \(j=0\), can be expressed in terms of the tangential components of the fields at the boundary

\[
R(\mathbf{H}', \mathbf{H}'') = \omega \int_{\partial \Gamma} (\mathbf{E}'' \times \mathbf{H}'') \cdot \mathbf{n} - (\mathbf{E}' \times \mathbf{H}') \cdot \mathbf{n}, \quad \text{when } j = 0.
\]

Consider a medium that is locally isotropic and such that \(\varepsilon(\mathbf{x})\) is purely imaginary (which may be a good approximation at low frequencies in conductive media) while \(\mathbf{\mu}(\mathbf{x})\) is purely real (and so (4.7) only holds in the limiting sense with \(\alpha_\delta = 0\)). We can write

\[
\varepsilon = i \mathbf{I} / (\omega \rho), \quad e = -i \omega \rho \mathbf{I} \quad \text{and} \quad \mathbf{\mu} = \mathbf{\mu} \mathbf{I},
\]

where the resistivity \(\rho(\mathbf{x})\) and the permeability \(\mathbf{\mu}(\mathbf{x})\) are purely real and positive. Then, the expression (4.36) reduces to

\[
R(\mathbf{H}', \mathbf{H}'') = \int \omega \rho (|j' - \nabla \times \mathbf{H}'|^2 - |j'' - \nabla \times \mathbf{H}''|^2) + 2 \omega^2 \mathbf{\mu} \mathbf{H}' \cdot \mathbf{H}'',
\]
and then, when \( j=0 \), the saddle-point variational principle (4.35) corresponds to the Borcea (1999) variational principle. (Note that \( j \) has a different meaning in that paper.) Now Maxwell’s equations still hold if we make the replacements

\[
\varepsilon \rightarrow -i\varepsilon, \quad \mu \rightarrow i\mu, \quad E \rightarrow iE, \quad B \rightarrow iB, \quad D \rightarrow D, \quad H \rightarrow H, \quad j \rightarrow j. \tag{4.41}
\]

(which is a special case of the transformation (2.57)), and with these replacements applied to (4.39), the expression (4.36) reduces to

\[
R(H', H'') = \int_{\Gamma} -2\omega\rho(j' - \nabla \times H') \cdot (j'' - \nabla \times H'') + \omega^2 \mu(|H'|^2 - |H''|^2), \tag{4.42}
\]

and, when \( j=0 \), the saddle-point variational principle (4.35) then corresponds to the other Borcea variational principle. Multiplying \( 3 \) and dividing \( m \) by \( e^{i\theta} \), where \( 0 > \theta > -\pi/2 \), leads to a continuous family of saddle-point variational principles that interpolate the two Borcea variational principles.

5. Minimization variational principles for electromagnetism

To obtain a minimization principle, we rewrite the constitutive law in the form (2.35), where \( \mathcal{L} \) is given by (2.36). Then, we introduce the quadratic functional

\[
Y(E', H'') = \int_{\Gamma} \left( \frac{\Omega E'}{-\mathcal{G}'} \right) \cdot \mathcal{L} \left( \frac{\Omega E'}{-\mathcal{G}'} \right) + 2\omega j' \cdot E', \tag{5.1}
\]

where \( \mathcal{G}' \) is given by

\[
\mathcal{G}' = \omega \begin{pmatrix} H'' \\ -j'' - \nabla \times H'' \end{pmatrix} \tag{5.2}
\]

(to ensure that \( -\omega j'' + \mathcal{G}' = 0 \) inside \( \Gamma \)), and \( E' \) and \( H'' \) have the same tangential components at the boundary \( \partial \Gamma \) as, respectively, \( E' \) and \( \omega H'' \). Given two vector fields \( s \) and \( t \), with zero tangential components at the boundary \( \partial \Gamma \), then (4.23) and the key property (4.19) imply

\[
\int_{\Gamma} \mathcal{G}' \cdot (\Omega s) + \omega j' \cdot s - \mathcal{T} \cdot (\Omega u'') = 0, \quad \text{where } \mathcal{T} = \omega \begin{pmatrix} t \\ -\nabla \times t \end{pmatrix}. \tag{5.3}
\]

As a result, we deduce that

\[
Y(E' + s, H'' + t) = Y(E', H'') + \int_{\Gamma} \left( \frac{\Omega s}{-\mathcal{T}} \right) \cdot \mathcal{L} \left( \frac{\Omega s}{-\mathcal{T}} \right). \tag{5.4}
\]

Consequently, because \( \mathcal{L} \) is positive definite under the assumption (4.7), we have the minimization variational principle

\[
Y(E', H'') = \inf_{E'} \inf_{H''} Y(E', H''), \tag{5.5}
\]

where the infimums are over vector fields \( E' \) and \( H'' \), with the same tangential components at the boundary \( \partial \Gamma \) as, respectively, \( E' \) and \( H'' \). Using (2.37), the
minimum value of $Y(E', H'')$ is

$$Y(E', H'') = \int_0^{\omega^2} (H' \cdot B'' - H'' \cdot B' + E' \cdot D'' - E'' \cdot D') + 2\omega j' \cdot E'. \quad (5.6)$$

Now the power dissipated into heat in the body $I$, averaged over a cycle of oscillation, is

$$W(E, H) = \frac{1}{2} \int_I \left[ \frac{\omega}{2\pi} \int_0^{2\pi} (E e^{-i\omega t})' \cdot \frac{\partial (D e^{-i\omega t})'}{\partial t} + (H e^{-i\omega t})' \cdot \frac{\partial (B e^{-i\omega t})'}{\partial t} 
+ (E e^{-i\omega t})' \cdot (j e^{-i\omega t})' \ dt \right]$$

$$= \frac{1}{2} \int_I \omega (H' \cdot B'' - H'' \cdot B' + E' \cdot D'' - E'' \cdot D') + E' \cdot j' + E'' \cdot j'', \quad (5.7)$$

which is proportional to (5.6) when $j=0$. Therefore, the minimum value of $Y$ has a physical interpretation when $j=0$.

Owing to the diagonal structure (4.20) of $Z$, the formula (5.1) reduces to

$$Y(E', H'') = \int_I \left( \frac{\omega}{2\pi} \int_0^{2\pi} \left( \nabla \times H'' - j'' \right) \cdot \mathcal{E} \left( \nabla \times H' - j' \right) 
+ \left( \nabla \times E' \right) \cdot \mathcal{M} \left( \nabla \times E' \right) \right) + 2\omega j' \cdot E', \quad (5.8)$$

where $\mathcal{E}(x)$ and $\mathcal{M}(x)$ are the positive definite matrices

$$\mathcal{E} = \begin{pmatrix} \epsilon'' + \epsilon' (\epsilon'')^{-1} \epsilon' & \epsilon'(\epsilon'')^{-1} \\ (\epsilon'')^{-1} \epsilon' & (\epsilon'')^{-1} \end{pmatrix},$$

$$\mathcal{M} = \begin{pmatrix} m'' + m'(m'')^{-1} m' & m'(m'')^{-1} \\ (m'')^{-1} m' & (m'')^{-1} \end{pmatrix}, \quad (5.9)$$

in which $m = -\mu^{-1}$. Also, the constitutive law (2.35) reduces to

$$\begin{pmatrix} D'' \\ E'' \end{pmatrix} = \mathcal{E} \begin{pmatrix} E' \\ -D' \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} H' \\ -B' \end{pmatrix} = \mathcal{M} \begin{pmatrix} B'' \\ H'' \end{pmatrix}. \quad (5.10)$$

Therefore, given any arbitrary prescribed values for the tangential components $E'$ and $H''$ at the boundary $\partial I$, the minimum in the variational principle (5.5) will satisfy the Euler–Lagrange equation and therefore will correspond to a solution of Maxwell equations. From the minimizing fields $E'$ and $H''$, one can, according to (5.10), determine

$$\begin{aligned}
E'' &= (\epsilon'')^{-1} \epsilon' E' + (\epsilon'')^{-1} (\nabla \times H'' - j'') / \omega, \\
H' &= m'(m'')^{-1} H'' - [m'' + m'(m'')^{-1} m'] (\nabla \times E') / \omega,
\end{aligned} \quad (5.11)$$

and, in particular, one can determine the tangential components of these fields at the boundary.

Proc. R. Soc. A (2009)
It is often the case that $\mu$ is real, while $\varepsilon$ has a positive definite imaginary part. In this event, we redefine $\mathbf{Z}$, $\mathbf{G}$ and $j$ by multiplying them by $e^{i\theta}$. Again, there is a continuous two-parameter family of related variational principles. We rewrite

Maxwell’s equations (4.1) in the form

$$
\nabla \times \tilde{\mathbf{E}} = i\omega \tilde{\mathbf{B}}, \quad \nabla \times \tilde{\mathbf{H}} = -i\omega \tilde{\mathbf{D}} + \tilde{j}, \quad \tilde{D} = \tilde{\varepsilon}\tilde{E} \quad \text{and} \quad \tilde{B} = \tilde{\mu}\tilde{H},
$$

(5.12)

where

$$
\begin{align*}
\tilde{E} &= e^{-i\omega t_0}E, \quad \tilde{D} = e^{i(\theta-\omega t_0)}D, \quad \tilde{H} = e^{i(\theta-\omega t_0)}H, \\
\tilde{B} &= e^{-i\omega t_0}B, \quad \tilde{j} = e^{i(\theta-\omega t_0)}j, \quad \tilde{\varepsilon} = e^{i\theta}\varepsilon \quad \text{and} \quad \tilde{\mu} = e^{-i\theta}\mu,
\end{align*}
$$

(5.13)

in which the real constant $t_0$ (which may be viewed as a time interval) can be chosen freely, while $\theta$ must be chosen, so that

$$
\tilde{\varepsilon}''(x) \geq \alpha_1 I \quad \text{and} \quad \tilde{\mu}''(x) \geq \alpha_2 I, \quad \forall x \in \Gamma,
$$

(5.14)

for some $\alpha_1 > 0$ and $\alpha_2 > 0$. Then, the fields with the tildes on them also satisfy Maxwell’s equations and the variational principles directly apply to them. With the subsequent replacement (5.13), we obtain a new set of variational principles parametrized by $\theta$ and $t_0$. When $\theta$ is non-zero, the minimum value of the minimization principles no longer has a physical interpretation as the time-averaged power dissipation in the body.

In particular, by choosing $\theta = 0$ and $t_0$, with $e^{-i\omega t_0} = -i$, we recover the dual variational principle that

$$
\tilde{Y}(E'', H') = \int_\Gamma \left( \frac{\omega E''}{\omega E''} \right) \cdot \mathcal{E} \left( \frac{\omega E''}{\omega E''} \right) \\
+ \left( \frac{\omega H'}{\omega H'} \right) \cdot \mathcal{M} \left( \frac{\omega H'}{\omega H'} \right) + 2\omega j'' \cdot E''
$$

(5.15)

is minimized at $E'' = E''$ and $H' = H'$, under the constraint that $E''$ and $H'$ have the same tangential values on $\partial\Gamma$ as, respectively, $E''$ and $H'$. A corollary is that

$$
W(E, H) \equiv |Y(E', H') + \tilde{Y}(E'', H')|/(4\omega)
$$

(5.16)

is minimized at $E = E' + iE'' = E$ and $H = H' + iH'' = H$, under the constraint that $E$ and $H$ have the same tangential values on $\partial\Gamma$ as, respectively, $E'$ and $H'$. By making the substitutions (5.13) in (5.6), it is easy to check that the minimum value of $W(E, H)$ is the average power dissipation $W(E, H)$ given by (5.7), even when $j$ is non-zero. An appealing feature of the expression (4.39) is that it remains invariant under the replacements $E \rightarrow e^{-i\omega t_0}E$ and $H \rightarrow e^{-i\omega t_0}H$ for any real value of $t_0$. However, a disadvantage is that the boundary conditions are overprescribed: if one imposes boundary conditions on the tangential values of $E$ and $H$ that are not compatible with the moduli $\varepsilon(x)$ and $\mu(x)$, then the fields that minimize $W(E, H)$ will not satisfy Maxwell’s equations.

When $\mu''$, or equivalently $m''$, is very small, we have

$$
\begin{align*}
\left( \nabla \times E' \right) \cdot \mathcal{M} \left( \nabla \times E' \right) &\approx \left[ m' \nabla \times E' - \omega H'' \right] \cdot (m'')^{-1} \left[ m' \nabla \times E' - \omega H'' \right],
\end{align*}
$$

(5.17)
and the variational principle is only useful if we take trial fields with $\mathbf{H}'' \approx m'\nabla \times \mathbf{E}'/\omega$. If we take $\mathbf{H}'' = m'\nabla \times \mathbf{E}'/\omega$, we still obtain a useful variational principle in the limit when $\mu$ is real. From (5.5) and (5.8), we have

$$Y(E') = \inf_{E'} Y(E''),$$

(5.18)

where

$$Y(E') = \int_{\Gamma} \left( \frac{\omega E'}{- \nabla \times \mu^{-1}(\nabla \times E')/\omega - j''} \right) \cdot \mathcal{E} \left( - \nabla \times \mu^{-1}(\nabla \times E')/\omega - j'' \right) + 2\omega j' \cdot E'.

(5.19)

In this variational principle, the infimum is over vector fields $E'$ with prescribed tangential components of $E'$ and $\mu^{-1} \nabla \times E'$ at the boundary $\partial \Gamma$. The infimum is attained at a field $E'$ associated with the solution of Maxwell’s equations satisfying the prescribed boundary conditions. When there are discontinuities in $\mu$ across some interior interfaces, the trial field $E'$ should be chosen so that the tangential component of $\mu^{-1}(\nabla \times E')$ is continuous across these interfaces.

6. Application to electromagnetic tomography

Here, we consider the tomography problem: how can one recover information about the functions $\varepsilon(x)$ and $\mu(x)$ given the measurements of the tangential components $E_{\|}(x)$ and $H_{\|}(x)$ of the fields $E(x)$ and $H(x)$ at the boundary $\partial \Gamma$? When it is known that the body $\Gamma$ consists of small inclusions in a matrix, this question has been answered by Ammari et al. (2001; see also the review by Ammari & Kang 2004). For other geometries, a different approach is needed and the variational principles provide this. First, observe from (4.15) and (5.7) that the power dissipation $W(E, H)$ depends only on the tangential values of these fields at the boundary

$$W(E, H) = -\frac{1}{2} \int_{\partial \Gamma} (E'_{\|} \times H'_{\|}) \cdot n + (E''_{\|} \times H''_{\|}) \cdot n,$$

(6.1)

and, as could be expected physically, is the time-averaged value of the flux of the Poynting vector into the body $\Gamma$. Thus, when $j=0$ and the trial fields $E'$ and $H''$ are chosen having the same tangential values at the boundary as, respectively, $E'$ and $H''$, the inequality

$$W(E, H) \leq Y(E', H'')/(2\omega)$$

(6.2)

provides, for a given $\varepsilon(x)$ and $\mu(x)$, a bound on the possible tangential components of the fields $E''$ and $H'$. Conversely, if these tangential components have been measured, the inequality provides one integral constraint on the possible values of $\varepsilon(x)$ and $\mu(x)$. Additional inequalities on $\varepsilon(x)$ and $\mu(x)$ may be obtained by choosing different trial fields, by working with other minimization variational principles in the family parametrized by $t_0$ and $\theta$, and by conducting a series of physical experiments with various different sets of values of the tangential components of the fields $E(x)$ and $H(x)$ at the boundary $\partial \Gamma$. It is not necessary that the moduli $\varepsilon(x)$ and $\mu(x)$ remain the same for each physical experiment, provided we have some model for how they vary from experiment to
experiment. For example, the frequency \( \omega \) could be different for the different experiments if we know that the permittivity has the form \( \varepsilon(x) = \varepsilon_0 + i\sigma_0/\omega \) (or some other known parametric dependence on \( \omega \)). The objective is then to constrain the possible values of the real functions \( \varepsilon_0(x) \) and \( \sigma_0(x) \).

One criticism of this approach is that it uses only information (in the case \( \theta = 0 \)) about the time-averaged power dissipation in the body \( \Gamma \). When \( t=\theta = 0 \), the measured values of the functions \( E'' \) and \( H' \) around the boundary \( \partial \Gamma \) are not used except to calculate the single scalar quantity \( W(E, H) \). This seems to be a tremendous waste of information. However, some extra information from the measured fields can be incorporated if one has measurements of the boundary value fields \( \mathbf{E}^{(j)}_\parallel(x) \) and \( \mathbf{H}^{(j)}_\parallel(x) \) for \( n \) experiments indexed by the integer \( j=1, 2, \ldots, n \), with the moduli \( \varepsilon(x) \) and \( \mu(x) \) remaining the same for each experiment. Then, one can infer the boundary values \( \mathbf{E}_\parallel \) and \( \mathbf{H}_\parallel \) associated with a linear combination of these fields,

\[
\mathbf{E}_\parallel = \sum_{j=1}^{n} \lambda_j \mathbf{E}^{(j)}_\parallel \quad \text{and} \quad \mathbf{H}_\parallel = \sum_{j=1}^{n} \lambda_j \mathbf{H}^{(j)}_\parallel ,
\]

where \( \lambda_1, \lambda_2, \ldots, \lambda_n \) are a set of complex constants. Letting \( \mathbf{E} \) and \( \mathbf{H} \) denote the corresponding electric and magnetic fields inside \( \Gamma \), it follows that \( W(E, H) \) can be expressed as the sum

\[
W(E, H) = \sum_{j=1}^{n} \sum_{k=1}^{n} (\lambda'_j \lambda'_k + \lambda''_j \lambda''_k) W_{jk} + (\lambda''_j \lambda'_k - \lambda''_k \lambda'_j) S_{jk}
\]

\[
= \frac{1}{2} \sum_{j=1}^{n} \sum_{k=1}^{n} (\lambda'_j \lambda'_k + \lambda''_j \lambda''_k) (W_{jk} + W_{kj} + (\lambda''_j \lambda'_k - \lambda''_k \lambda'_j) (S_{jk} - S_{kj}),
\]

involving the real quantities

\[
W_{jk} = -\frac{1}{2} \int_{\partial \Gamma} \left( \mathbf{E}'^{(j)}_\parallel \times \mathbf{H}'^{(k)}_\parallel \right) \cdot \mathbf{n} + \left( \mathbf{E}''^{(j)}_\parallel \times \mathbf{H}''^{(k)}_\parallel \right) \cdot \mathbf{n},
\]

\[
S_{jk} = -\frac{1}{2} \int_{\partial \Gamma} \left( \mathbf{E}'^{(j)}_\parallel \times \mathbf{H}''^{(k)}_\parallel \right) \cdot \mathbf{n} - \left( \mathbf{E}''^{(j)}_\parallel \times \mathbf{H}'^{(k)}_\parallel \right) \cdot \mathbf{n},
\]

which can be extracted from the experimental measurements of the boundary fields. Thus, \( W(E, H) \) depends not only on the \( n \) diagonal elements, \( W_{jj} \), which physically represent the time-averaged power dissipation in each experiment, but also on the \( n(n-1) \) elements \( W_{jk} + W_{kj} \) and \( S_{jk} - S_{kj} \), with \( j > k \), which have no such physical interpretation. It makes sense to choose trial fields that are also linear combinations,

\[
\mathbf{E}' = \sum_{j=1}^{n} (\lambda_j \mathbf{E}^{(j)}_\parallel)' = \sum_{j=1}^{n} \lambda'_j \mathbf{E}'^{(j)} - \lambda''_j \mathbf{E}''^{(j)},
\]

\[
\mathbf{H}'' = \sum_{j=1}^{n} (\lambda_j \mathbf{H}^{(j)}_\parallel )'' = \sum_{j=1}^{n} \lambda_j \mathbf{H}''^{(j)} + \lambda_j \mathbf{H}'^{(j)}.
\]

Then, \( Y(E', H'')/(2\omega) - W(E, H) \) must be a positive quadratic form in the \( 2n \) real variables \( \lambda'_1, \lambda''_1, \lambda'_2, \lambda''_2, \ldots, \lambda'_n, \lambda''_n \), and the positive semi-definiteness of the associated \( 2n \times 2n \) matrix will give constraints on the functions \( \varepsilon(x) \) and \( \mu(x) \).
We can also use the variational principle (5.18) for tomography purposes. When \( \mu \) is real and \( j=0 \), the inequality
\[
W(E, H) \leq Y(E^/) / (2\omega)
\] (6.7)
applies. Often one knows in advance that the material inside \( \Gamma \) is non-magnetic, so that \( \mu = I \) inside \( \Gamma \). Then, from knowledge of \( W(E, H) \), (6.7) provides a constraint on the possible values of the electrical permittivity \( \varepsilon \) for every suitable choice of the trial field \( E^/ \). On the other hand, suppose that the medium inside \( \Gamma \) was two phase, with an unknown phase boundary and with two different values of \( \mu \) inside \( \Gamma \). Then, the variational principle (6.7) will be useless unless the choice of the trial field is carefully correlated with the unknown phase boundary, so that the tangential component of \( \mu^{-1}(\nabla \times E^/) \) is continuous across the phase boundary.

When \( j\neq0 \), then we have the inequality
\[
W(E, H) \leq W(E, H),
\] (6.8)
where \( W(E, H) \) is given by (5.16) and the trial fields \( E \) and \( H \) are required to have the same tangential values as, respectively, \( E \) and \( H \). Given the measurements of the tangential values of \( E \) and \( H \), and the trial fields \( E \) and \( H \), we expect that this inequality should provide a useful constraint on the possible values of \( \varepsilon(x) \) and \( \mu(x) \).

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