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Nanoporous carbon/WO₃ anodes for an enhanced water photooxidation

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1. Introduction

Advances on the use of sunlight as a sustainable low-cost source of energy in various fields have rather been limited to the use of inorganic semiconductors - TiO₂ and metal oxides and sulfides as CdS, Fe₂O₃, or ZnO among most representatives-, that should be stable, nontoxic, inexpensive, and have suitable electronic band positions for an efficient visible light absorption [1–5]. WO₃ is a n-type semiconductor usually presenting a 3D arrangement of slightly distorted corner-shared [WO₆] octahedra compared to the ideal cubic perovskite (type ReO₃), responsible of its electrooptical, electrochromic, ferroelectric and catalytic properties [6–10]. It is considered an interesting visible-light stable photocatalytic material due to its relatively large abundance, non-toxicity, physical and chemical resilience in harsh environments, and most importantly its strong absorption features in the solar spectrum (i.e., band gap between 2.4 and 2.8 eV). The valence band edge of WO₃ is located at approximately 2.9 V vs Ag/AgCl (at pH 0) and is positive enough for the oxidation of water (i.e., high oxidation power of the photogenerated holes). However, the potential associated with the conduction band edge is not negative enough (ca. 0.3 V vs Ag/AgCl) for the reduction of water without a bias voltage. This provokes the surface accumulation of electrons, increasing the recombination rate and ultimately reducing the photoelectrochemical performance of the photoanode [9–13].

To compensate for this limitation of WO₃, several strategies have been adopted such as surface sensitization, optimization of the particle size and morphology control, incorporation of additives and transition/noble metal dopants, or the application in association with suitable photocathodes -i.e. p-type semiconductors-in tandem photoelectrochemical cells [4,14–20]. All these approaches aim to increase the electron trapping by inhibiting the charge recombination.

In this study we have investigated the use of nanoporous carbon additives to WO₃ as an alternative to improve the performance of...
the semiconductor towards the photooxidation of water. Based on the high electronic density of carbon materials, provided by the stacked graphene-layers, it seems reasonable to anticipate that the above-mentioned surface accumulation of electrons can be dissipated by delocalization in the π-electron density of the graphitic structure of carbons. Indeed, the use of carbons of varied nature and characteristics (e.g., carbon nanotubes and nanofibers, graphene, nanoporous carbons, graphitic carbon nitrides) as additives to semiconductors has been widely explored ([21–27] and references therein), and several studies report the enhanced photocatalytic yields of such semiconductor/carbon catalysts. In most cases the improved performance is attributed to various phenomena such as: i) an enhanced visible light absorption of the hybrid materials; ii) synergistic effects based on nanocorrelation on the porosity of the carbon support, iii) the strong carbon/semiconductor interfacial electronic effects that favor the reaction with electron donor/accep-

orators present in the medium. More recently, the photochemical activity of nanoporous carbons has been reported based on the efficient conversion of light into chemical reactions in the confinement of the nanopore space, offering interesting perspectives in different fields [28–35].

Considering all the above, the objective of this work was to explore the effect of a nanoporous carbon as additive to WO₃ in the performance of the resulting hybrid WO₃/NC catalysts for the photooxidation of water. Thus, we provide a systematic investigation of the photoelectrochemical responses of WO₃/nanoporous carbon electrodes for the oxidation of water as a function of the irradiation wavelength, the applied potential bias and the amount of carbon additive. The mixtures were prepared by adding increasing amounts of a nanoporous carbon (ranging from 5 to 80 wt%) to WO₃ powders; the carbon additive was chosen based on its photochemical activity under varied illumination conditions towards the photooxidation of phenol in aqueous solution [28,31].

The novelty of the work is to provide an insight on the role of the carbon additive in the external quantum yield for water oxidation as a function of the carbon ratio, and to correlate the photoelectrochemical response of the WO₃/carbon anodes with the wavelength of the irradiation source. Despite the modest incident photon-to-current conversion efficiency values obtained, our results demonstrate the higher photocatalytic conversion of the photonic solar after the incorporation of the nanoporous carbon, due to the enhanced electron collection provided by the carbon additive. The photoelectrochemical response of the tungsten oxide nanoparticles (severely limited by a poor carrier mobility of the semiconductor) was enhanced by the incorporation of low amounts of the carbon additive; this offers a new approach for the processing of photoelectrocatalysts using low cost nanoporous carbon as sustainable metal-free additives to semiconductors.

2. Experimental

2.1. Materials

The carbon additive (sample referred to as NC) used in the preparation of the WO₃/carbon mixtures is an amorphous nanoporous carbon. Details on the origin and characteristics of this material have been presented in detail elsewhere [26,28]. We herein re introduce some textural and chemical properties for a better comprehension of the response of the WO₃/carbon photoanodes towards the photooxidation of water. Commercially available WO₃ nanopowders (Fluka, puriss 99.9%, average particle size ca. 40–60 nm, surface pH 4.3) with a monoclinic crystalline structure (JCPDS 01-083-0950), were used as semiconductor.

2.2. Preparation of the electrodes

For the fabrication of the electrodes, about 100 mg of solids (either bare WO₃ or WO₃/carbon mixtures) were dispersed in 750 μL of isopropyl alcohol; about 100 μL of this ink were spread by a micropipette onto an FTO-coated glass slide (surface resistivity ~ 7 Ω/sq) and spin-coated at 3500 rpm for 30 s. The films were dried at 60 °C for 30 min and then annealed at 300 °C for 1 h (heating rate 5 °C/min). The annealing temperature was carefully chosen to avoid modifications in the carbon matrix and semiconductor structure. Indeed, neither the crystallinity of the semiconductor nor the surface chemistry and textural properties of the carbon additive were modified upon the annealing treatment applied in the preparation of the electrodes (see Figs. S1–S3 in the Electronic Supporting Information File, ESI). The WO₃/carbon electrodes will be labeled as W/NCₓ, where x stands for the amount of nanoporous carbon additive. A series of electrodes containing ca. 5, 10, 20, 50 and 80 wt% of carbon were prepared.

2.3. Photoelectrochemical measurements

Photoelectrochemical measurements were carried out using an Ivium CompactStat potentiostat/galvanostat. The studies were performed in a standard three-electrode system using an electrochemical cell with an optically flat circular quartz window (diameter 2 cm). A Pt wire and Ag/AgCl (saturated KCl) were used as counter and reference electrode, respectively. The prepared electrodes were immersed in an aqueous solution containing 0.5 M Na₂SO₄ (inert electrolyte), previous adjustment of the solution pH to 1.3 by the addition of H₂SO₄. All the experiments were carried out at room temperature. Transient photocurrent responses under on/off illumination cycles were obtained at a constant bias potential between −0.1 and +0.6 V vs. Ag/AgCl. Current equilibrium at each applied potential was allowed before the irradiation (dark). The electrochemical behavior of the electrodes was also explored by cyclic voltammetry at a potential sweep of 20 mV s⁻¹. The photoelectrochemical responses were measured several times in different electrodes to evaluate the reproducibility; average values corresponding to at least 3 distinct electrodes are presented in this work. Measurements under illumination were performed using a blue LED (Thorlabs) featuring a narrow emission spectrum centered at 371 nm. The LED was powered by a waveform generator (Stanford Research Systems), with a photon flux of 3.4 × 10¹⁵ cm⁻² s⁻¹. Incident-photon-to-current efficiency (IPCE) at varied wavelengths (ranging between 360 and 460 nm) was measured using 150 W Arc Xe lamp fitted to a single-pass monochromator. The intensity of the irradiation sources was calibrated using a standard silicon photodiode (Newport Corporation). The IPCE values were calculated as:

\[
\text{IPCE}(\%) = \frac{1240 J_{ph}}{\lambda P} \times 100
\]

where \(\lambda\) (nm) is the wavelength of the monochromatic light irradiating the electrode, \(J_{ph}\) (A/cm²) corresponds to the stationary photocurrent density measured under illumination with P (W/cm²) incident irradiation. The number 1240 carries the matching units.

2.4. Characterization techniques

The porosity of all the samples was determined by N₂ adsorption/desorption isotherms at −196 °C in a volumetric analyzer (Micromeritics). Before the experiments, the samples were outgassed at 120 °C for 17 h under vacuum (10⁻⁴ Torr). The specific surface area (S_{BET}), total pore volume (V_{POR}), and micropore volume (Wo N₂) were obtained from the gas adsorption isotherms.
Micropore volumes were evaluated using the Dubinin-Radushkevich equation. UV–Vis diffuse reflectance spectra were recorded on a Shimadzu spectrometer (UV-2501) equipped with an integrating sphere and using BaSO₄ as a blank reference.

2.5. Electrochemical impedance spectroscopy measurements

Electrochemical impedance spectroscopy (EIS) measurements were carried out in a 3-electrode cell in 0.5 M Na₂SO₄, using WO₃ as working electrode, and a Pt wire and Ag/AgCl (saturated KCl) as counter and reference electrodes, respectively. Potentiostatic EIS measurements were recorded in a potentiostat/galvanostat (VMP-3, Biologic) in the frequency range of 0.01 Hz–100 kHz for selected potentials between −0.6 and +0.6 V vs. the Ag/AgCl and using 10 mV AC amplitude. Data obtained was fitted to the Mott–Schottky (MS) equation (see below) that shows the dependence of the interfacial capacitance (C) with the applied potential (E) under depletion conditions (assuming that the dimensions of the semiconductor crystal are larger than the width of the space charge layer), providing useful information about the changes in the charge distribution in the passive layer at the electrode/electrolyte interface.

\[
\frac{1}{C_S^2} = \frac{1}{C_H^2} + \frac{2}{\varepsilon \varepsilon_0 N} \left( E - E_{fb} - kT/e \right)
\]

where \( C_S \) is the charge capacity, \( C_H \) the Helmholtz capacity, \( N \) the carrier density (donor \( N_D \) or acceptor \( N_A \)), \( \varepsilon \) the dielectric constant of the electrode material, \( \varepsilon_0 \) the electric permittivity of vacuum, \( e \) the elementary charge, \( k \) the Boltzmann constant, \( T \) the absolute temperature and \( E_{fb} \) the flat band potential.

### Table 1

|                  | \( S_{BET} \) [m² g⁻¹] | \( V_{MES} \) [cm³ g⁻¹] | \( W_0 N_2 \) [cm³ g⁻¹] |
|------------------|-------------------------|-------------------------|-------------------------|
| WO₃              | 6                       | 0.038                   | 0.002                   |
| W/NC5            | 56                      | 0.077                   | 0.019                   |
| W/NC10           | 108                     | 0.104                   | 0.037                   |
| W/NC20           | 205                     | 0.160                   | 0.070                   |
| W/NC50           | 524                     | 0.315                   | 0.187                   |
| W/NC80           | 839                     | 0.472                   | 0.279                   |
| NC               | 1033                    | 0.520                   | 0.320                   |

a. Total pore volume evaluated at p/p₀ ~0.99.
b. Micropore volume evaluated from the Dubinin-Radushkevich method.

### 3. Results and discussion

Fig. 1 shows the voltammetric response of the WO₃/carbon electrodes under dark and square-wave light perturbation centered at 371 nm. For WO₃, the voltammograms displayed the characteristic shape of an n-type semiconductor with two distinctive regions (accumulation and depletion) depending on the bias potential. The cathodic current recorded when the electrode potential is swept below 0.1 V vs. Ag/AgCl is attributed to the reduction of the W⁶⁺ ionic state due to the intercalation of protons and the electrochemical injection of the excess electrons from the substrate to the conduction band of the semiconductor (according to \( WO_3 + xH^+ + xe^- \rightarrow H_xWO_3 \)) [36,37]. This charge transfer reaction provokes a blue coloration in the electrode (known electrochromism behavior of WO₃). Upon the reverse scan, WO₃ recovers
the translucent-greenish state as a result of the reversible deintercalation of the protons and the simultaneous oxidation of tungsten. Above 0.1 V, the voltammogram exhibits a flat current corresponding to the depletion region where electrons can no longer flow. By expanding the potential window above 1.8 V, an anodic faradaic current appears due to electrochemical decomposition the electrolyte (water oxidation, results not shown).

For the WO3-carbon electrodes, the capacitive contribution of the nanoporous carbon due to the formation of the electrical double layer is clearly seen in the voltammograms. The current densities under dark conditions gradually increase with the amount of carbon of the electrodes, in agreement with the increasing porous features of the WO3/carbon mixtures (Table 1, Figs. S4–S6).

In the case of the electrode with the highest amount of carbon (i.e., 80 wt%) it becomes difficult to distinguish the shape associated to the redox reactions at the semiconductor, due to the large capacitive response of the carbon (Fig. S4). Additionally, for this sample, an oxidation/reduction reversible hump is observed close to 0.3 V vs. Ag/AgCl. This is attributed to the pseudo-faradic contribution of the surface groups of the carbon matrix (mainly quinone/hydroquinone pair) [38,39], as the similar behavior is observed in the voltammogram of the nanoporous carbon alone (Fig. S6). Furthermore, the contribution of the double layer capacitance due to the nanoporous carbon allowed verifying the real catalysts composition. Experimental capacitance values measured in all the WO3/NC electrodes matched the theoretical values corresponding to the amount of carbon additive incorporated to the mixtures (Fig. S5), confirming the homogenous dispersion of the catalysts inks during the preparation of the anodes for all the carbon ratio.

Upon squared-waved illumination at 371 nm, anodic photocurrents were recorded in the voltammetry profiles of all the electrodes, with the exception of W/NC80 (Fig. 1). The electrons in the valence band of WO3 are promoted to the conduction band and transferred through the bulk semiconductor to the back contact; meanwhile, the holes are transported to the electrode/electrolyte interface where they may react with electron donors present in the
The photocurrent onset potential is associated to the conduction band and the optical band gap obtained by diffuse reflectance spectroscopy (Fig. S7). In this regard, the spectrum of bare WO₃ presented the characteristic absorption sharp edge above 470 nm, corresponding to a band gap value of 2.6 eV, in agreement with those previously reported for this material [8–14]. The characteristic absorption sharp edge of WO₃ was still observed in the diffuse reflectance spectra of the WO₃/carbon mixtures (Fig. S7); the characteristic fingerprint of these spectra suffered a gradual increase in the absorbance in the region between 500 and 800 nm (with the amount of carbon added), due to the strong light absorption features of the carbon component. A similar behavior has been reported for other semiconductor/carbon mixtures (based on TiO₂ and Bi₂WO₆ semiconductors and using carbon nanotubes, activated carbons and carbon hydrochars as additives) [23,25,41–43].

To further evaluate the photoelectrochemical response of the semiconductor we have performed impedance measurements under dark conditions (Mott-Schottky representation is shown in Fig. S8). The dependence of the capacitance values of the space charge region yielded a straight line with a positive slope, characteristic of n-type semiconductors [40] that is inversely proportional to the donor density (N_D). The corresponding value obtained for WO₃ was N_D = 1.9 × 10²⁰ cm⁻³, which is in good agreement with those reported in the literature for monoclinic tungsten trioxide [14,44]. In the case of the WO₃/carbon mixtures the situation is more complex as the electrochemical impedance is dominated by the large double layer capacitance of the nanoporous carbon (as also seen in Fig. 1), and there is a large incertitude in the values of the dielectric constant of the mixtures.

Fig. 2 shows examples of the photocurrent transients at various potentials for WO₃ and WO₃/NC10 electrodes (data corresponding to all the studied electrodes is shown in Fig. S9). Photocurrent transients were registered for all studied electrodes at different potentials to discriminate the photoelectrochemical response from the capacitive contribution, particularly in the electrodes with the largest amounts of carbon additive.

In all cases, the transient photocurrent curves upon on/off illumination exhibited the characteristic shape of an n-type semiconductor. On switching-on the light, an initial rapid and sharp increase was observed, followed by a gradual decay as the photogenerated electrons and holes recombine. The current retracted to its original value once the illumination was switched-off instantaneously and it was reproducible during repeated on/off illumination cycles for various electrodes demonstrating the reproducibility and stability of the electrodes under successive cycles; only for higher applied potentials (0.4–0.6 V vs Ag/AgCl) and for electrodes with small carbon loading (ca. 5–20 wt%) a slow decay is observed which we attribute to the stabilization of the electrons in intermediate surface states. The photocurrent onset potential is quite similar in all cases, with a significant increase in photocurrent magnitude in the presence of the carbon additive. It can also be observed that, besides the increase in photocurrent values, the main features on the photocurrent transients remain unaffected by the addition of carbon. On the other hand, the sharp potential dependence of the photocurrent magnitude suggests that the driving force for electron transport is strongly dependent on the position of the electrode Fermi energy. It should be clarified that electron transport is mainly determined by diffusion, given that the nanostructured nature of the oxide prevents the formation of depletion layers [2]. For some electrodes, an initial anodic current peak (overshoot) is observed (Fig. 2, Fig. S9), which is associated to fast recombination processes and complex carrier trapping effects discussed further below.

Fig. 3A shows a direct comparison of the photocurrent transients obtained at 0.6 V vs Ag/AgCl for all the electrodes, to allow a better observation of the effect of the incorporation of the carbon additive. As seen, the photocurrent responses show a sharp increase upon addition of small quantities of carbon, showing the larger value close to 20 wt% for the electrodes with carbon content of 50 wt% and below; the recorded photocurrent values are higher than those of the bare semiconductor. This is most remarkable considering that the amount of semiconductor is gradually reduces as the amount of carbon additive increases, demonstrating that
these W/NC electrodes are more photoactive than bare WO$_3$. The lower photocurrents of the electrodes with the highest carbon content is most likely due to the lower photogeneration of carriers in the electrodes showing a lower content of WO$_3$ and to shadowing effects from the carbon additive. Indeed, if the photocurrent responses are normalized by the WO$_3$ content (Figs. S10 and S11), a plateau for carbon amounts of 20–50 wt% is observed. It is remarkable that even the normalized photocurrent of the electrode W/NC80 is higher than that of bare WO$_3$ electrode.

The photocurrent decay observed in the transients (Figs. 2 and 3A) originates from a complex relationship between potential distribution across the electrode, trapping of electrons or holes in intermediate surface states and recombination losses (arising from accumulation of holes and electrons at the surface and in the bulk, respectively) [45]. Recombination losses upon accumulation of holes and electrons can be discarded since both phenomena typically manifest by a significantly sharper photocurrent decay accompanied by an equally large cathodic current (undershoot) when the light is turned off (and electrodes in the conduction band react with the accumulated holes). Furthermore, similar photocurrents were obtained by front- and back-side illumination modes of the electrodes, discarding carrier transport limitations as the main mechanism governing the photocurrent losses [46,47]. Hence, the photocurrent decay is most likely attributed to the trapping of electrons in deep states, which leads to a redistribution of the interfacial potential. This process can be regarded as a shift of the effective potential towards more a negative value, promoting a decay of the photocurrent.

Fig. 3B shows the photocurrent transients recorded under consecutive illumination cycles at different time scales. Upon initial illumination, photogenerated electrons are captured by deep trap states, leading to the interfacial potential redistribution mentioned above (Fig. 3B). If a subsequent illumination cycle is allowed before the surface states are emptied, the transient displayed almost a square-shape reaching a stationary photocurrent value. However, if the electrode is kept at open-circuit conditions for 5 min in the dark, the photocurrent transients showed again the gradual fall in the initial photocurrent associated to various trapping states. Similar changes in the population of deep trap states in nanostructured semiconductors in the second time-scale have been

Fig. 5. SEM and TEM images of the some selected photoanodes (first raw) and the catalyst mixtures illustrating the dispersion of the semiconductor and carbon components in the inks.
largely depends on the formation of aggregates of both component. Interaction between both phases is rather weak (dispersive) and seem the optimal method for the preparation of catalysts, since the electrodes were reproducible. Composition of the inks was homogeneous (Fig. S5) and hence the TEM images in Fig. 5 and Fig. S12. As above-mentioned, the conductor and carbon particles in the inks is shown in the SEM and further improvement. In this regard, the dispersion of the semi-conductors has not yet been optimized, thus there should be room for carbons).

ration of the photoanodes (metal-free and amorphous nanoporous considering the nature of the carbon additive used in the preparation. In other words, the fall in the carrier generation due to the diffusion length of carriers brought about by the carbon additive. Also, IPCE values showed a strong dependence with the bias potential for all the electrodes (Fig. 4B).

Although these IPCE values are not as high as those reported for other water splitting photocatalysts [47], they are most outstanding concerning the nature of the carbon additive used in the preparation of the photoanodes (metal-free and amorphous nanoporous carbons).

It should also be pointed out that the preparation of the electrodes has not yet been optimized, thus there should be room for further improvement. In this regard, the dispersion of the semiconductor and carbon particles in the inks is shown in the SEM and TEM images in Fig. 5 and Fig. S12. As above-mentioned, the composition of the inks was homogeneous (Fig. S5) and hence the electrodes were reproducible.

However, the physical mixing of the two components would not seem the optimal method for the preparation of catalysts, since the interaction between both phases is rather weak (dispersive) and largely depends on the formation of aggregates of both component. An improved close contact of the two materials would be expected to facilitates the charge diffusion through the mixture, and hence improving the performance of the electrodes.

To investigate the wavelength dependence of the photochemical response of the electrodes, we recorded the photoelectrochemical response of the anodes under different illumination conditions. Fig. 6A shows the spectral response of the photocurrent for the W/NC electrodes between 360 and 460 nm. In the case of WO3, the maximum conversion was recorded at wavelength about 370–400 nm, in agreement with data reported in the literature [49].

The profile of the W/NC photoanodes is very similar to that of bare WO3. Our previous studies have shown that this nanoporous carbon is photoactive towards the oxidation of phenol in a different spectral range, showing a poor conversion efficiencies at 400 and 450 nm [29], which suggests that different mechanisms apply for these reactions (phenol and water photooxidation). As already observed at 371 nm (Fig. 3A), higher photocurrents were obtained for all the W/NC electrodes, confirming that the effect of the carbon additive applies for a wide wavelength range (and not exclusively in the ultraviolet irradiation). The effect of the carbon matrix is more clearly seen in Fig. 6B, where the photocurrent densities of the W/NC electrodes (j_w/nc/wo3) are normalized vs that of the bare WO3 (j_WO3). For W/NC10 and W/NC20 the higher values of this ratio are due to the higher efficiency in the photon to current conversions. The profiles show a somewhat constant dependence with the wavelength, with photocurrent densities for W/NC electrodes was found between 3 and 4 times higher than for bare semiconductor for all wavelengths studied (Fig. 6B). The slight increase in the j_w/nc/wo3 ratio over 400 nm, although not very pronounced is significant, and points out that the effect of the carbon additive is higher at wavelengths corresponding to the visible range of the spectrum.

Summarizing, enhanced photocurrent and IPCE values were recorded for the W/NC photoanodes, confirming a more efficient charge carrier collection when the nanoporous carbon additive is incorporated to the electrodes. To understand this behavior we must consider several factors. First of all, the smaller diffusion

![Fig. 6. Dependence of the (A) IPCE values and (B) normalized photocurrent densities (j_w/nc/wo3) with the wavelength of the illumination source (bias potential of 0.6 V, irradiation with a Xe lamp fitted to a single-pass monochromator). (A color version of this figure can be viewed online.)](image-url)
length of the minority carriers in nanoporous electrodes (due to the nanometric dimensions of the pore voids) is expected to reduce their recombination, as the carriers easily reach the surface and react with electron acceptors. Secondly, the conjugated $\pi$-electron density of the carbon material could act as an acceptor of the photogenerated electrons upon illumination of WO$_3$, rapidly boosting the electron diffusion via delocalization within the graphene-like sheets. An increased electron diffusion length has already been reported for other semiconductors/carbon materials based on carbon nanostructures with high charge carrier mobility (such as nanotubes and graphenes). For instance, Yousefzadeh and col. [50], described the enhancement in photocurrent density by increasing the MWCNT content from 0 to 1 wt% in the WO$_3$ films associated to the high electron conductivity, charge transfer and long electron lifetime; more recently Fu and co-workers [51] reported the use of WO$_3$-reduced graphene oxide nanocomposite showing an enhancement in the photocatalytic performance due to the highly facilitated electron transfer. We herein demonstrate that this behavior can be achieved with a better yield using low-cost nanoporous carbons with a disordered turbostratic structure. In this regard, we have recently reported the visible light driven photoelectrochemical water photooxidation splitting on metal-free nanoporous carbons with varied surface functionalization [30,52,53]. The ability of nanoporous carbons with amorphous structure to oxidize water has been linked to the presence of chromophoric groups and the generation of carriers in reactive sites (associated to surface groups or free edges sites). Moreover, the beneficial effect of carbon additives of varied nature on the photoactivity of WO$_3$ nanopowders has also been demonstrated for other reactions (i.e., photooxidation of pollutants in aqueous solution) [43,54].

4. Conclusions

WO$_3$/nanoporous carbon photoanodes were explored to increase the light conversion efficiency of the semiconductor in the photooxidation of water within a wide spectral range. Higher conversion efficiencies in the range of 360–460 nm were obtained upon the incorporation of small amounts of a nanoporous carbon additive, compared to the performance of the bare semiconductor. This is attributed to the improved electron collection provided by the carbon component, which lowers recombination losses, thereby facilitating the oxidation of water at the WO$_3$/nanoporous carbon/electrolyte interface. Higher IPCE values were obtained after the incorporation of the carbon additive to the semiconductor; nevertheless, the modest IPCE values and the strong potential dependence of the photocurrent indicate that the recombination and charge collection are still limiting the efficiency. Hence, the possibility of achieving higher conversion efficiencies upon illumination in a wide spectral range using nanoporous carbons with different porous features and conductivity, and the effect of the nanopore confinement on the light propagation for the photoelectrochemical oxidation of water at semiconductor/nanoporous carbon anodes is currently under investigation.

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Appendix A. Supplementary data

Supplementary data related to this article can be found at http://dx.doi.org/10.1016/j.carbon.2016.07.045.

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