Strain-controlled synthesis of ultrathin hexagonal GaTe/MoS\textsubscript{2} heterostructure for sensitive photodetection

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Highlights
High-quality ultrathin hexagonal GaTe is synthesized via strain engineering

Phase transition mechanism of GaTe is revealed by using theoretical calculation

Excellent photosensing properties are observed in GaTe/MoS\textsubscript{2} p-n heterostructure

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Strain-controlled synthesis of ultrathin hexagonal GaTe/MoS2 heterostructure for sensitive photodetection

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SUMMARY

Ultrathin hexagonal GaTe, with relatively high charge density, holds great potential in the field of optoelectronic devices. However, the thermodynamical stability limits it fabrications as well as applications. Here, by introducing two-dimensional MoS2 as the substrate, we successfully realized the phase-controlled synthesis of ultrathin h-GaTe, leading to high-quality h-GaTe/MoS2 heterostructures. Theoretical calculation studies reveal that GaTe with hexagonal phase is more thermodynamically stable on MoS2 templates, which can be attributed to the strain stretching and the formation energy reduction. Based on the achieved p-n heterostructures, optoelectronic devices are designed and probed, where remarkable photoresponsivity (32.5 A/W) and fast photoresponse speed (<50 μs) are obtained, indicating well-behaved photo-sensing behaviors. The study here could offer a good reference for the controlled growth of the relevant materials, and the achieved heterostructure will find promising applications in future integrated electronic and optoelectronic devices and systems.

INTRODUCTION

Ultrathin GaTe semiconductors have attracted considerable attention due to their wonderful electrical and optoelectronic properties and are expected to be promising materials for high-sensitive photodetectors (Zolyomi et al., 2013; Lin et al., 2018; Yuan et al., 2015; Wang et al., 2015a, 2015b). Most III-VI layered compound, such as GaS, InSe, and GaSe, shares a hexagonal phase which is thermodynamically stable and can be easily achieved with ultrathin thickness (Han et al., 2020; Yu et al., 2019; Li et al., 2020a, 2020b; Mudd et al., 2013; Cai et al., 2018). As for GaTe, it possesses a different thermodynamically stable phase, namely monoclinic layered structure, with C2/m space group symmetry, which is very difficult to be obtained with ultrathin thickness (Cai et al., 2017; Dong et al., 2019; Cui et al., 2009). From this point of view, the experimental fabrication of hexagonal GaTe (h-GaTe) is meaningful for the decreasing of the dimension of GaTe. Previous studies have indicated that the formation energy of h-GaTe is higher than that of monoclinic-GaTe (m-GaTe), leading to that h-GaTe is a metastable phase and is hardly to be achieved (Yu et al., 2019; Cai et al., 2017). Up to now, most experimental studies on GaTe and its heterojunction are based on the monoclinic structure rather than the hexagonal phase (Wang et al., 2014, 2015a, 2015b; Liu et al., 2014; Yang et al., 2016; Yan et al., 2018; Gang et al., 2014; Fonseca et al., 2016). Unlike the phase transition of transition-metal chalcogenides (Zhang et al., 2017; Song et al., 2016; Sung et al., 2017; Cho et al., 2015; Liu et al., 2018; Yin et al., 2017; Zhou et al., 2016; Yoo et al., 2017), there is no effective strategy in realizing the phase transition of GaTe from m-phase to h-phase. It is highly desired to develop a controllable preparation method for making GaTe with high-purity hexagonal phase and to realize a high-performance photodiode device by designing the heterostructure based on ultrathin GaTe.

Here, we demonstrate, for the first time, the controllable growth of high-quality ultrathin h-GaTe on MoS2 templates. Theoretical studies further reveal that h-GaTe is more thermodynamically stable on MoS2 than m-GaTe, which can be attributed to the strain stretching and formation energy reduction of h-GaTe. Microscopic and spectroscopic characterizations are employed to characterize the achieved h-GaTe/MoS2 vertical heterostructures, indicating high crystal quality. Moreover, excellent photovoltaic and photosensing properties are observed in the h-GaTe/MoS2 vdW heterostructure. Under illumination, photo-excited electron-hole pairs are readily separated by the large built-in electrical field formed at the GaTe-MoS2...
interface, leading to fast response speed (<50 μs) and high photoresponsivity (32.5 A/W). The study here could offer a good reference for the controlled growth of the relevant materials, and the obtained ultrathin GaTe/MoS2 heterostructures will provide potential materials for application in integrated electronic and optoelectronic devices and systems.

RESULTS AND DISCUSSION

Synthesis of ultrathin h-GaTe on MoS2 templates

Figure 1A shows the schematic diagram of the synthesis of h-GaTe/MoS2 vertical heterostructure. During the growth process of h-GaTe/MoS2 heterostructure, MoS2 was first synthesized on the SiO2/Si substrate via the typical chemical vapor deposition (CVD) growth strategy, and then, h-GaTe was fabricated on the as-achieved MoS2 template. The monolayer GaTe is the configuration of Te–Ga–Ga–Te repeating units (in Figure S1). The detailed descriptions of the growth procedures are presented in the experimental method. Figure 1B exhibits the large-scale microscope image of the achieved samples, where partially covered ultrathin h-GaTe/MoS2 vertical heterostructures are uniformly distributed on the substrate. Density-functional theory (DFT) calculations were further conducted to reveal the growth thermodynamics in h-GaTe/MoS2 heterostructures (Sánchez-Royo et al., 2002; Leao and Lordi, 2011; Rahmilow et al., 2007; Kresse and Furthmüller, 1996; Blöchl, 1994; Kresse and Joubert, 1999; Klimes et al., 2011). The relative stability between the h-phase and the m-phase is related to the formation energy difference between them. Considering the two phases have different the total numbers of atoms in their unit cells, we define $\Delta E = E_h - E_m$, where $E_h$ and $E_m$ represent the total energy of the h-phase and m-phase per GaTe unit, respectively.

$\Delta E$ is about 6 meV for the free-standing condition, that is, the h-GaTe monolayer is less stable than the m-GaTe monolayer in this case, which is in accordance with the case for the bulk GaTe. On the contrary, while MoS2 is introduced as the substrate, $\Delta E$ is reduced to −12 meV, suggesting that the h-phase is more energetically favorable than the m-phase. Generally, a substrate may impose strain and form bonding to the overlayer. However, in our model, given that both the overlayer and substrate are semiconducting and the interactions between them are vdW forces, the bonding effect is expected to be weak, and the strain should be the main reason that changes the $\Delta E$. The interface structure for the h-phase GaTe consists of a $\sqrt{7} \times \sqrt{7}$ supercell of GaTe and a $2\sqrt{3} \times 2\sqrt{3}$ supercell of MoS2, which gives rise to a lattice mismatch of about 0.9% between the two constituents. For the m-phase, a rectangular supercell that contains sixteen unit cells of the monoclinic GaTe monolayer and ninety unit cells of the MoS2 monolayer is built for GaTe/MoS2, where the lattice mismatch of 0.5% is measured. It means that the MoS2 can definitely induce strain in the upper GaTe and influence its stability, leading to the synthesis of h-GaTe. Table S1 gives the calculated formation energy of $E_h$ and $E_m$ for both freestanding and supported cases.
Optical characterization of ultrathin h-GaTe/MoS2 vertical heterostructures

Raman and photoluminescence (PL) spectra were performed to investigate the chemical composition distributions and the local optical properties of the GaTe/MoS2 heterostructures. Figure 2A shows a typical as-grown GaTe/MoS2 sample with partially covered structures, where the components can be clearly identified from the optical contrast. Then, we performed thickness identification of the selected red area shown in Figure 2A, and the corresponding Atomic Force Microscope (AFM) image is depicted in Figure 2B. It indicates that the thickness of the top GaTe is about 0.8 nm, indicating the monolayer nature of h-GaTe. Figure 2C shows the spatial-resolved Raman intensity mapping at 130 cm\(^{-1}\) of the sample in Figure 2A, where strong Raman signal is collected at the lower part of the sample, corresponding to the A\(_{1g}\) mode of h-GaTe. We also collected the Raman spectra at the lower part of the sample, and the result is shown in Figure 2E, where the Raman peaks of both MoS\(_2\) and GaTe are clearly observed, indicating the formation of h-GaTe/MoS\(_2\) heterostructure (the monoclinic version of the Raman scattering as comparison in Figure S2 of the supplemental information). Room temperature PL mapping (670 nm–690 nm) and PL spectra are further collected and shown in Figures 2D and 2F, respectively. The weak PL emission signal from the section covered by GaTe is shown in Figure S3 of the supplemental information. It is clear that the uncovered MoS\(_2\) region gives a strong PL emission, but a very weak PL signal can be collected in the MoS\(_2\) region covered with GaTe. Such a typical PL quenching phenomenon observed in GaTe/MoS\(_2\) can be attributed to the type-II band alignment at the interface which can efficiently separate the photoexcited electron-hole pairs (Figure 2G) (Li et al., 2018, 2019; Yang et al., 2017). All the results demonstrate that the as-grown heterostructures are of high quality without alloying.

TEM images of h-GaTe/MoS\(_2\) heterostructure

We further used transmission electron microscopy (TEM) combined with energy-dispersive spectroscopy (EDS) to analyze the microstructure and chemical composition of the bilayer vertical h-GaTe/MoS\(_2\) heterostructure. Figure 3A demonstrates the low-magnification TEM image of an as-transferred GaTe/MoS\(_2\) heterostructure on Cu grid, where the heterointerface can be obviously identified from the contrast. Figure 3B
shows the distribution of Ga, Te, Mo, and S in the yellow rectangle region in Figure 3A. It is obvious that Mo and S are evenly distributed throughout the whole area, while Ga and Te can only be found in the triangle region. The EDS measurement is utilized to further confirm the chemical composition of the as-grown MoS2/GaTe heterostructure. Figure 3C displays the EDS spectra recorded from two different positions marked in the heterostructure. The element ratio of Ga to Te is further confirmed to be 49:51, matching well with the GaTe stoichiometric ratio. The chemical states of Ga and Te in the h-GaTe were characterized by X-ray photoelectron spectra where binding energy peaks of Ga 3d3/2, Te3 d3/2, and Te3 d5/2 are observed at 25.5 eV, 583 eV, and 54.7 eV, respectively, indicating the existence of GaTe (Figure S4). Two sets of hexagonal symmetrical diffraction spots in the same direction can be clearly identified in the selected area electron diffraction (SAED) graph in Figure 3D, corresponding to the MoS2 lattice (2.7 Å) and h-GaTe lattice (3.5 Å), respectively. On the contrast, SAED pattern of m-GaTe is a parallelogram (Figure S5), which is totally different from that of h-GaTe. The high-resolution TEM image of the overlapped region in Figure 3E displays apparent Moire patterns, which is attributed to the overlapping lattices between GaTe and MoS2, matching well with the atomic structure model of the GaTe/MoS2 heterostructure in Figure 3F. All the above characterizations clearly illustrate that the as-fabricated vertical h-GaTe/MoS2 heterostructures are of good crystal quality.

Non-linear optical properties of GaTe flakes

To study the stacking mode of h-GaTe layers, we further explored second harmonic generation (SHG) performance of h-GaTe and determined the symmetry of the sample. SHG is a nonlinear optical phenomenon, which is ultrsensitive to the centrosymmetry of 2D materials. It is noteworthy that the MoS2 obtained in our experiment possesses 2H phase with a centrosymmetry structure, leading to weak SHG signals in bilayer...
2H-MoS2 (Figure S6). Here, a GaTe/MoS2 heterostructure integrated with bilayer MoS2 is chosen to investigate the SHG performance of the GaTe, and the SHG characteristics of the samples were collected with transmission geometry. As shown in Figure 4A, an expected single peak located at 400 nm is observed, which is exactly at the doubled frequency of the femtosecond (fs)-laser power of 800-nm excitation wavelength. Moreover, the SHG intensity monotonically increases with the enhancement of the fs-laser power. The detailed relationship between SHG intensity and laser power is plotted in Figure 4B, in which a slope of 2.05 can be acquired, in accordance with the theoretical value of 2, demonstrating the second-order nonlinear optical process (Zeng et al., 2019; Saynatjoki et al., 2017). We also investigate the relationship between the thickness and the SHG of the GaTe in Figure 4C. As can be seen, the SHG intensity is enhancing with the increase of sample thickness, suggesting the h-GaTe flakes are 3R stacking mode, with the atomic structure shown in Figure 4D. The results indicate that the achieved sample can be potentially used in nonlinear optics.

Photoresponse characterizations of bilayer GaTe/MoS2 heterostructures
To further investigate the electrical charge transport properties and optoelectronic performance of the resulted ultrathin GaTe/MoS2 vertical heterostructure, the optoelectronic devices based on the obtained bilayer GaTe/MoS2 p-n heterojunctions were investigated schematically under vacuum at room temperature. Figure 5A exhibits the configuration of the h-GaTe/MoS2 photodetector device, where Cr/Au electrodes are separately fabricated on monolayer MoS2 and heterostructure region. The transport characteristics in Figure 5B demonstrate an obvious current rectification behavior with $V_g$ from 0 to 60 V, which is common with the behaviors that are observed in traditional p-n junctions (electrodes 2 and 3). The inset image of Figure 5B displays a typical optical image of the as-fabricated bilayer p-n heterojunction device (device area: 42 µm$^2$). Figure S7 also shows the output characteristics under measured using difference electrodes contact from 0 to 60 V. Figure 5C plots the $I_{ds}$-$V_{ds}$ curves under 520-nm light irradiation with power densities changing from 0.01 to 78.5 mW/cm$^2$ at zero gate bias. Obviously, the photocurrent increases gradually with the increase of light power intensity. From the magnified $I_{sc}$-$V_{ds}$ curves shown in Figure 5D, apparent photovoltaic behaviors are observed with maximum open-circuit voltage reaching up to 0.145 V at a power density of 78.5 mW/cm$^2$. Figure 5E displays the extracted short-circuit current ($I_{sc}$) and open-circuit voltage ($V_{oc}$) as a function of incident light power density. It is obvious that both $I_{sc}$ and $V_{oc}$

![Figure 4. Non-linear optical properties of GaTe flakes](image-url)
increase gradually with the increase of the incident light power density. Figure 5F plots the curves as a function of the $V_{ds}$ under different light power density. The generated electrical power ($P_{EL}$) in the p-n diode can further be calculated with equation $P_{EL} = I_{ds} \times V_{ds}$. It suggests that the maximum $P_{EL}$ is generated at 0.075 V with a light intensity of 78.5 mW/cm$^2$. It also means that the most suitable output working voltage is 0.075 V for a photovoltaic device. The photocurrent at $V_{ds} = 1$ V is extracted and plotted in Figure 5G, where the photocurrent ($I_{ph}$) is defined as $I_{ph} = I_{light} - I_{dark}$. The relationship between photocurrent and laser power density can be fitted within an equation of $I_{ph} = aP^\alpha$, where $a$ is the parameter related to the photodetector responsivity, $P$ is the laser power density, and $\alpha$ is dimensionless exponent of the power law. The parameters $a$ and $\alpha$ are determined to be 3.2 and 0.51, respectively. The responsivity ($R$) can be calculated with the equation of $R = I_{ph}/PA$, where $A$ is the effective area of the device channel. Benefiting from the ultrathin structure and the large built-in electrical field at the junction interface, the photo-generated electron-hole pairs can be effectively separated, resulting in high responsivity (Li et al., 2017). The maximum photoresponsivity is measured to be 32.5 A/W at an incident power of 0.6 mW/cm$^2$, which is higher than...
that of previous reports on multilayer m-GaTe/MoS2 vertical heterostructures.\cite{13-15, 19} Time-resolved photoresponse measurements were performed to further investigate the photoresponse speed of the h-GaTe/MoS2 heterojunction photodetector, and the results are shown in Figure S5H. As can be seen, by periodically turning on and off the laser (520 nm, 30.4 mW/cm²), the photocurrent can be switched well between the on and off state, showing excellent stability and reliability. Further investigation indicates that the device possesses rise time ($t_{\text{rise}}$) of about 25 μs and fall time ($t_{\text{decay}}$) of about 48 μs, which is faster than the pure GaTe and GaTe/MoS2 vertical heterostructure-based device (Figure S5) \cite{Wang et al., 2014, 2015a, 2015b}. Table S2 gives the comparison of the photoresponse performance of our h-GaTe/MoS2 p-n vertical junction with several other GaTe reports. Photoswitching properties of the device illuminated by different wavelengths (520 nm and 800 nm) are further carried out in Figure S8, where stable photoresponse at each wavelength is observed. Although the photocurrent illuminated by the 800-nm laser is weak than the 520-nm laser, it is surprising that the device can respond not only to visible light but also to near-infrared (NIR) light. The light response to NIR light may be originated from the interlayer excitation, which is beyond both the MoS2 and the GaTe absorption edge.

Conclusions

In summary, we demonstrate the successful synthesis of high-quality ultrathin h-GaTe/MoS2 heterostructures. Theoretical calculation studies reveal that the achieved h-GaTe is more stable on 2D MoS2 substrate than m-GaTe, which leads to greater strain stretching and reduces the formation energy of h-GaTe. Excellent photovoltaic and photosensing properties are observed in the ultrathin p-n GaTe/MoS2 heterojunctions. With the drain bias voltage fixed at 1 V, high photoresponsivity of 32.5 A/W has been realized at light of 520 nm. Furthermore, ultra-fast response times $t_{\text{rise}}$ (25 μs) and $t_{\text{decay}}$ (48 μs) are also proven. The study here could offer a good reference for the controlled growth of the relevant materials, and the achieved GaTe/MoS2 heterostructure will find promising applications in future integrated electronic and optoelectronic devices and systems.

Limitations of the study

In addition to the ultrathin h-GaTe being prepared on MoS2 templates in this study, more relevant phase transition materials can be prepared and applied for more breakthroughs in structure control and performance enhancement. Considering the limitation of the experiment platform, the detailed process of in situ phase transition remains unknown.

STAR METHODS

Detailed methods are provided in the online version of this paper and include the following:

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SUPPLEMENTAL INFORMATION

Supplemental information can be found online at https://doi.org/10.1016/j.isci.2021.103031.

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AUTHOR CONTRIBUTIONS

F.L., X.Z., D.L., and A.P designed the research idea and experiments. M.C. carried out DFT calculation. Y.W., X.Z., Z.Z., D.Z., J.Y., and Z.L. analyzed the date. All authors discussed the results and commented on the manuscript.

DECLARATION OF INTERESTS

The authors declare no competing financial interest.

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STAR METHODS

KEY RESOURCES TABLE

| REAGENT or RESOURCE | SOURCE | IDENTIFIER |
|---------------------|--------|------------|
| Chemicals, peptides, and recombinant proteins | | |
| Gallium(III) iodide | Alfa Aesar | CAS: 13450-91-4 |
| Tellurium powder | Alfa Aesar | CAS: 13494-80-9 |
| Sulfur powder | Alfa Aesar | CAS: 7704-34-9 |
| Molybdenum(VI) oxide | Alfa Aesar | CAS: 1313-27-5 |
| Sodium chloride | Alfa Aesar | CAS: 7647-14-5 |
| Software and algorithms | | |
| Origin 8.1 | OriginLab | https://www.originlab.com/ |

RESOURCE AVAILABILITY

Lead contact
Further information and requests for resources and materials should be directed to and will be fulfilled by the Lead Contact, Prof. Anlian Pan (anlian.pan@hnu.edu.cn).

Materials availability
This study did not generate new unique reagents.

Date and code availability
- Data reported in this paper will be shared by the lead contact upon request.
- This paper does not report original code.
- Any additional information required to reanalyze the data reported in this paper is available from the lead contact upon request.

METHOD DETAILS

Synthesis of the 2D h-GaTe/MoS₂ heterostructure
The GaTe/MoS₂ heterostructures were synthesized by a two-step growth method using CVD technology. The first-step growth of monolayer MoS₂ is a traditional process. The mixed powder of 0.5mg sodium chloride (NaCl) and 2 mg Molybdenum(VI) oxide (MoO₃) was selected as the solid source for the first-step growth. For vertical GaTe/MoS₂ heterostructures, as-grown MoS₂ templates were placed at the downstream of a CVD system. The temperature was gradually heated up to 690°C in 20 min, and then kept the temperature for 5 min. Meanwhile, the 5mg Te powder was put in the upstream region about 450°C and the 2mg GaI₃ powder was put in the low temperature region about 210°C. A mixture flow of 100 sccm Ar/H₂ (with 10% H₂) was operated as the carrier gas. After the growth process, the furnace was cooled down to room temperature.

Characterization
The optical images of heterostructures were characterized using optical microscopy (Zeiss Axio Scope A1). The optical measurements (PL and Raman) were performed with the confocal-PL system (WiTec, alpha-300). A 520 nm solid state laser (power: about 5 mW, spot size: 1–2 μm) was used to excite the samples. For atomic-structural characterizations measurements, The STEM and EDS measurements were carried out on JEOL ARM200F microscope operated at 200 kV and equipped with a probe-forming aberration corrector. Standard e-beam lithography (EBL, Raith 150) and metal thermal evaporation were performed to fabricate the Au/Cr (10 nm/50 nm) electrodes on the produced heterostructures. The electrical and optoelectronic properties of the heterostructures were measured in vacuum with Lake Shore Probe Station and Agilent B1500A semiconductor analyzer at room temperature.
Density functional theory calculations
DFT calculations were performed using the Vienna Ab Initio Simulation Package. The pseudopotentials were constructed by the projector augmented wave method. Van der Waals dispersion (vdW) forces between the GaTe monolayer and MoS$_2$ substrate were accounted for through the opt PBE-vdW functional by using the vdW-DF method. A 12 × 12 Monkhorst-Pack k-mesh was used to sample the surface Brillouin zone (BZ) for the hexagonal phase and a 6 × 6 k-mesh was used for the monoclinic phase. For the interface structure, i.e., GaTe/MoS$_2$, a 6 × 6 was used for the hexagonal phase (Table S1). While for the monoclinic phase, a 1 × 3 k-mesh was used. A plane-wave energy cutoff of 300 eV was used for structural relaxation and total energy calculations.

QUANTIFICATION AND STATISTICAL ANALYSIS
Each value of SHG intensity shown in Figure 4B of the main manuscript corresponds to the average value obtained from multiple measurements at a given excitation power.