Multiswitching Synchronization of Non-Identical Hyperchaotic Lorenz and Chen Systems

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Abstract: In this paper, synchronization of four-dimensional hyperchaotic Lorenz system, (drive system) with the Chen’s system (response system) are investigated based on backstepping technique. Unlike the well-known i = j, multiswitching of the indices was employed in the usual master-slave synchronization scheme. We provided varieties of non-identical indices, that is, i = j of the driver system. In this high dimension, more switching options for constructing the error space vector due to the large number of variables are available; and by implication providing variety of synchronization directions between variables of the master and the slave systems, which could be used in securing electronic information and communication.

Keywords: multiswitching, synchronization, hyperchaotic systems

1. Introduction

It is generally known that nonlinear deterministic dynamical systems are ubiquitous in nature and that they are well known for their unusual property of sensitive dependence on initial conditions which gives rise to their temporal complexity and apparent randomness. This property indeed leads to several intriguing attributes that have been investigated for many decades in the fields of sciences, medicine and engineering [1] as well as the study of nonlinear systems in mathematics [2]. In particular, chaotic dynamics, bifurcation structures, multistability of attractors, pattern formations, synchronization dynamics and many more have been rigorously investigated in attempts to elucidate and describe nonlinear behaviours and their potential applications.

One of the most important attributes of nonlinear dynamical systems is synchronization, originally proposed by Pecora and Carroll [3]. This concept has been considered as a major breakthrough in chaotic dynamics [4], due to its potential applications in modelling brain activities, chemical reactions; and more importantly, in information processing and secure communication [5]. These foreseen applications has triggered the enormous research attention given to chaos synchronization for over two decades. Increasing interest in the study of synchronization of chaotic systems has led to the discovery of various methods and types of synchronization.

In classifying chaos synchronization, we specifically mention here the original notion of identical/complete synchronization [3,6], general chaos synchronization [7], phase synchronization [8], lag synchronization [9], generalized synchronization [10,11], measure synchronization [12–14], projective synchronization [15–17], anticipated synchronization [18, 19], reduced-order synchronization [20], finite-time synchronization [21], and function projective synchronization [22].

A good number of previous works on chaos synchronization were aimed at achieving synchronization between state variables of the master systems and that of similar slave system. Notably, these conditions would undoubtedly constrain the synchronization to simultaneously take place in predictable and predetermined directions. In 2008, Ucar et al. [23] proposed the multiswitching synchronization of chaotic systems, in which the slave system’s state variables synchronizes with different state variables of the driver. Despite its clear relevance to information security, only a few studies have been dedicated to this kind of synchronization to-date [4] [24] – [29].

In addition, most multiswitching synchronization scheme have been devoted to lower dimensional chaotic systems (See [4] and references therein). A few have considered hyperchaotic systems which indeed show complex non-periodic behavior with at least two Lyapunov exponents (See Ref. [29] and references therein). However, it has been shown that secure communications based on low dimensional system does not guarantee absolute security due to the possibility of reconstructing the temporal complexity and apparent randomness; and thereby separating the signals from the secure information [6]. For this reason, the primary motivation of synchronization in secure communication which is to hide certain information to be transmitted in chaotic signal and retrieving by the technique of chaotic synchronization; maybe defeated by using low dimensional chaotic systems.

Motivated by the aforementioned, we examine in this paper, the multiswitching synchronization of a four-dimensional (4D) hyperchaotic Lorenz and Chen systems – a higher dimensional hyperchaotic system with at least two positive Lyapunov exponent and having the advantages of more switching options for constructing the error space vector due to the large number of available variables and more temporal complexity, with enriched dynamics providing varieties of synchronization directions between variables of the master and that of the slave systems. With such a
scheme, it is hoped that improved security would be achieved when applied in communication.

The rest of the paper is organized as follows: in section 2, we give a brief description of the 4D hyperchaotic Lorenz and Chen systems. Section 3, general definitions and formulations of concepts. We designed the controllers based on backstepping section 4. In section 5 we gave numerical simulation results and the paper is concluded in section 6.

2. System Description

2.1 Description of the Hyperchaotic Lorenz’s model

Based on the Lorenz system, a new chaotic system was reported by [31] which has attracted the attention of many researchers [32-35]. The Lorenz chaotic system is described by

\[
\begin{align*}
\dot{x} & = \alpha (y - x) \\
\dot{y} & = cx - xz - w \\
\dot{z} & = xy - bz \\
\dot{w} & = \delta yz
\end{align*}
\]

where \(a\), \(b\) and \(c\) are positive real constants. It is a 3D autonomous system with six terms including only two quadratic terms in a form very similar to the Chen and Lu systems in [36-39], but it has three very different fixed points: one saddle and two stable node-foci. By introducing a linear feedback controller to the second equation of the Lorenz like system (1), the following new hyperchaotic system is obtained [40]:

\[
\begin{align*}
\dot{x} & = \alpha (y - x) \\
\dot{y} & = cx + y - xz - w \\
\dot{z} & = xy - bz \\
\dot{w} & = \delta yz
\end{align*}
\]

For the purpose of this work, we let \(a=10, b=8/3, c=28, d=0.1\). The attractors of the system at the states \(x-y, x-z, x-w, y-z, y-w, z-x, z-y\) and \(z-w\) are as shown in the upper part of Figure 1 while the time series of the combined state of the system is as shown in the lower part.

2.2 Description of the Hyperchaotic Chen’s model

The hyperchaotic Chen system is given by

\[
\begin{align*}
\dot{x} & = \alpha (y - x) + w \\
\dot{y} & = dx - xz + cy \\
\dot{z} & = xy - bz \\
\dot{w} & = \delta yz + \epsilon w
\end{align*}
\]

Where \(a=35, b=2, c=2, d=1, e=0.5\) are parameters. The attractors of the system at the states \(x-y, x-z, x-w, y-z, y-w, z-x, z-y\) and \(z-w\) are as shown in the upper part while the time series is shown in the lower part of Figure 2.

3. Definition and Formulation of

4. Design of Controllers

Multiswitching Synchronization

Let us consider the following master-slave n dimensional chaotic systems, where the master systems is given by

\[
\begin{align*}
\dot{x}_1 & = f_1(x_1) \\
\dot{x}_2 & = f_2(x_2) \\
\dot{x}_m & = f_m(x_m)
\end{align*}
\]

and the controlled slave system is given by

\[
\begin{align*}
\dot{y}_1 & = g_1(y_1) \\
\dot{y}_2 & = g_2(y_2) \\
\dot{y}_n & = g_n(y_n)
\end{align*}
\]

where \(x_1, y_1 (i = 1, 2, ..., n) \in \mathbb{R}^n\) are state space vectors of the systems, \(f_m, g_n : \mathbb{R}^n \rightarrow \mathbb{R}^n\) are two continuous vector functions and \(u_i, (i = 1, 2, ..., n) : \mathbb{R}^n \rightarrow \mathbb{R}\) are nonlinear control functions.

Definition 1

According to [41] in [4], If there exists two constants matrices \(A, B \in \mathbb{R}^n\) and \(A, B = 0\), such that \(\lim_{k \to \infty} kBy = 0\) and where \(k.k.k\) is the matrix norm and \(A, B\) are scaling matrices, then systems (4) and (5) are said to be in synchronization.

Comment 1

The error states in relation to the definition 1 are strictly chosen to satisfy the definition \(e_i (i = j, 1, 2, 3...n)\) are the indices of the error and \(n\) refers to the number of dimension of the chaotic system. Hence, we can write e

\[
\begin{align*}
e_{ij} &= y_i - x_j
\end{align*}
\]

where \(e_{ij}\) is the error dynamical system is interchanged and \(\lim_{k \to \infty} kBy = 0\) then, systems (4) and (5) are said to be in complete synchronization, in this case there is no switching.

If the error states in relation to definition 1 are redefined such that \(e_i (i = j)\) or \(e_i (i = j)\) and any of the member in the error dynamical system is interchanged and \(\lim_{k \to \infty} kBy = 0\) then, systems (4) and (5) are said to be in multiswitching synchronization state.

Comment 2

If \(i\) and \(j\) are chosen such that for \(e_i = y_i - x_j\), \((i = j)\), then we have

\[
\begin{align*}
e_{11} &= y_1 - x_1, e_{12} &= y_1 - x_2, e_{13} &= y_1 - x_3, \ldots e_{1n} &= y_1 - x_n, \\
e_{21} &= y_2 - x_1, e_{22} &= y_2 - x_2, e_{23} &= y_2 - x_3, \ldots e_{2n} &= y_2 - x_n, \\
e_{31} &= y_3 - x_1, e_{32} &= y_3 - x_2, e_{33} &= y_3 - x_3, \ldots e_{3n} &= y_3 - x_n, \\
e_{41} &= y_4 - x_1, e_{42} &= y_4 - x_2, e_{43} &= y_4 - x_3, \ldots e_{4n} &= y_4 - x_n, \\
e_{mn} &= y_m - x_n
\end{align*}
\]

for \(m\) dimensions of the slave systems and \(n\) dimensions of the master systems. We comment here also that there are various switchings that can be formulated in addition to the above.

From system (4), we let \(x = x_1, y = y_1, z = z_1\) and \(w = w_1\) for the master system and \(x = x_2, y = y_2, z = z_2\) and \(w = w_2\) for the slave system, and we chose the following switchings
Let the master system be
defines the error between $e$, choosing $e_1 = e_{13}$, $z_1 = x_1$, $z_2 = x_2$, and $w_1 = y_1$.

and the slave system as

\begin{align}
\dot{x}_i &= a_i(x_i - x_l) \\
\dot{y}_i &= c_i x_i + y_l - x_l z_l - w_i
\end{align}

and substituting
\begin{align}
\dot{e}_{13} &= -a_2 e_{13} + a_2 (y_l - z_l) - x_1 y_1 + b_1 z_1 + w_1 + u_1 \\
\dot{e}_{14} &= z_1 - w_1 \\
\dot{e}_{15} &= w_2 - x_l
\end{align}

We stabilize (15) by describing the second Lyapunov function as
\[
v_2 = v_1 + \frac{1}{2} \eta_2
\]

whose time derivative is
\[
u_2 = -2 c_2 e_{22} + z_2 (c_1z_1 + 1) + y_1 (c_2 - 1) - z_1 (d_2 + x_l) - c_1 x_1 - w_1 - c_2 e_{22} + u_2
\]

By substituting for $\eta_1$ in (16), and choosing $e_{13} = a_2 (\eta_2)$ as a virtual controller, we can write
\[
\dot{\eta}_1 = v_1 + \frac{1}{2} \eta_2
\]

So that
\[
\dot{\eta}_1 = - (a_2 - k) \eta_1^3 - (c_2 - k) \eta_2^3
\]

which is $\leq 0$ for $a_2 > 0$, $c_2 > 0$, $k \leq 0$. Thus the subsystem is negative definite and asymptotically stable. Since the error between $e_{13}$ and $a_2 (\eta_2)$ is estimative as $\eta_1 e_{13} u_1 (\eta_1)$ and $a_2 (\eta_2) = 0$, let $e_{34}$ be $\eta_1$, we can write the ($\eta_1$, $\eta_2$, $\eta_3$) subsystem as
\[
\begin{align}
\eta_1 &= -(a_2 - k) \eta_1 + a_2 \eta_2 \\
\eta_2 &= -(c_2 - k) \eta_1 + b_2 \eta_3 \\
\eta_3 &= -b_2 e_{34} + e_{22} (x_2 + 1) + x_1 y_1 - b_2 w_1 - d_1 y_1 z_1 + u_3
\end{align}
\]

We can stabilize (20) by defining the third Lyapunov function as
\[
\dot{\eta}_3 = - (a_1 - k) \eta_1^2 - (c_1 - k) \eta_2^2 - (b_1 - k) \eta_3^2
\]

which is $\leq 0$ for $a_1 > 0$, $c_1 > 0$, $b_1 > 0$, $k \leq 0$. Thus the ($\eta_1$,$\eta_2$,$\eta_3$) subsystem is asymptotically stable. Since the error between $e_{13}$ and $a_2 (\eta_2)$ is estimative as $\eta_1 e_{22}$ - $a_2 (\eta_1)$ and $a_3 (\eta_1) = 0$, let $e_{41} = \eta_1$, we can write the ($\eta_1$, $\eta_2$, $\eta_3$, $\eta_4$) whole system as

\[
\dot{\eta}_4 = -c_4 e_{41} + e_{22} (x_2 + 1) + x_1 y_1 - b_2 w_1 - d_1 y_1 z_1 + u_3
\]

With error dynamics presented above, if appropriate $u_1$, $u_2$, $u_3$ and $u_4$ are chosen such that the system is stable and unchanged, then asymptotic stabilization would be realized leading to synchronization of the system. If $e_{13} = \eta_1$, its time derivative is $\dot{\eta}_1 = \dot{\epsilon}_{13}$ and we can stabilize the first part of (9), using the Lyapunov function, as
\[
v_1 = \frac{1}{2} \eta_1^2
\]

whose time derivative is
\[
\dot{v}_1 = \dot{\eta}_1^2
\]

By substituting for $\eta_1$ from (9) we have
\[
\dot{v}_1 = -(a_2 - k) \eta_1^2 + a_2 (y_l - z_l) - x_1 y_1 + b_1 z_1 + w_1 + u_1
\]

choosing $e_{13} = a_2 (\eta_2)$ as a virtual controller, we have
\[
u_1 = -a_2 (y_l - z_l) + x_1 y_1 + b_1 z_1 + k \eta_1
\]

So that
\[
\eta_1 = -(a_2 - k) \eta_1^2
\]

which is $\leq 0$ for $a_2 > 0$, $k \leq 0$. Thus the subsystem is negative definite and asymptotically stable. Since the error between $e_{13}$ and $a_2 (\eta_2)$ is estimative as $\eta_1 e_{13} u_1 (\eta_1)$ and $a_2 (\eta_2) = 0$, let $e_{34}$ be $\eta_1$, we can write the ($\eta_1$, $\eta_2$, $\eta_3$) subsystem as

\[
\eta_1 = -(a_2 - k) \eta_1 + a_2 \eta_2 \\
\eta_2 = -(c_2 - k) \eta_1 + b_2 \eta_3 \\
\eta_3 = -b_2 e_{34} + e_{22} (x_2 + 1) + x_1 y_1 - b_2 w_1 - d_1 y_1 z_1 + u_3
\]

We can stabilize (20) by defining the third Lyapunov function as
\[
\dot{\eta}_3 = - (a_1 - k) \eta_1^2 - (c_1 - k) \eta_2^2 - (b_1 - k) \eta_3^2
\]

which is $\leq 0$ for $a_1 > 0$, $c_1 > 0$, $b_1 > 0$, $k \leq 0$. Thus the ($\eta_1$,$\eta_2$,$\eta_3$) subsystem is asymptotically stable. Since the error between $e_{13}$ and $a_2 (\eta_2)$ is estimative as $\eta_1 e_{22}$ - $a_2 (\eta_1)$ and $a_3 (\eta_1) = 0$, let $e_{41} = \eta_1$, we can write the ($\eta_1$, $\eta_2$, $\eta_3$, $\eta_4$) whole system as

\[
\dot{\eta}_4 = -c_4 e_{41} + e_{22} (x_2 + 1) + x_1 y_1 - b_2 w_1 - d_1 y_1 z_1 + u_3
\]
We can stabilize (25) by defining the fourth Lyapunov function as
\[ V_4 = v_3 + \frac{1}{2} \eta_4^2 \] whose time derivative is
\[ \dot{V}_4 = \dot{v}_3 + \eta_4 \dot{\eta}_4 \]

We have
\[ \dot{v}_3 = -(c_2 - k) \eta_1^2 - (c_2 - k) \eta_2^2 - (b_2 - k) \eta_3^2 - (e - k) \eta_4^2 \]
which is ≤ 0 for \( a_2 > 0, b_2 > 0, c_2 > 0 \) and \( k ≤ 0 \) Thus, the whole system \((\eta_1, \eta_2, \eta_3, \eta_4)\) is asymptotically stable. For other cases (2-5), the controllers are as presented in equations (30), (31), (32) and (33) respectively.

5. Numerical Simulations

Here we present our numerical simulation in order to verify the effectiveness of controllers derived above and those in (30)-(33). We used the Runge-Kutta simulation tool. In all the figures presented, the hyperchaotic Lorenz system (drive) parameters selected remain constant at \( a_1 = 10, b_1 = \frac{8}{3}, c_1 = 28, d_1 = 2.25 \), with initial condition values \( x_1 = 0.1, y_1 = 0.1, z_1 = 0.1, w_1 = 0.1 \), the hyperchaotic Chen system (response), parameters selected remain constant at \( a_2 = 35, b_2 = 3, c_2 = 12, d_2 = 7, e = 0.5 \) and variable parameters chosen as \( x_2 = 0.5, y_2 = 0.01, z_2 = 0.8, w_2 = 0.5 \). The step size was maintained at \( H = 0.005 \) and \( t = 10 \). The synchronization of the slave system with the master system at each of the state variable is presented in (Figures 3) for states \( e_1, e_2, e_3 \) and \( e_4 \) respectively, while the fifth graph in this figure shows the result for the combined states for case 1. In each of these figures, synchronization took place when each of the controllers \( u_1, u_2, u_3 \) and \( u_4 \) was activated at \( t ≥ 10 \). Using the same parameters, and making use of our results for other cases 2, 3, 4 and 5 for the non-identical hyperchaotic systems unified and periodically converged to zero as shown in the figures 4 as time tends to infinity which signifies that the global synchronization between systems (6) and (7) has been achieved.

6. Conclusion

Summarily in this paper, we have proposed, analysed and validated the possibility of the multiswitching synchronization of the non-identical hyperchaotic Lorenz (drive) and that of Chen’s (response) systems based on integrator backstepping technique. We extended the usual master-slave synchronization scheme for low order chaotic systems to study the synchronization of this higher order systems on one side, and provided various switches in the...
design of the controllers. Each of the 4-dynamical states was successfully synchronized. The synchronization of each of the switches in other cases were also successful. By implication, electronic information can be hidden in any or all of this 4D hyperchaotic system and such information can be locked up in any of the states in each of the cases. Such information can be transferred, communicated and retrieved by applying the control inputs for each or all the dynamical states and respective switches. This makes the information more secure not only because of the hyperchaos status of the system in consideration, but also the several switches that must be unlocked to retrieve the information. Our numerical results confirm the effectiveness of the analytical technique and we believe that they are observable in laboratory experiments.

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Figure 1: Lorenz system. The upper part of the figure shows some of attractors of the system for planes x−y, x−z, x−w, y−z, y−w, z−x, z−y and z−w where synchronization could take place, while the lower part shows the time series of the whole system at $a_1=10$, $b=8/3$, $c=28$, $d_1=0.1$, are shown respectively.
Figure 2: Chen system. The upper part of the figure shows some of attractors of the system for planes x-y, x-z, x-w, y-z, y-w, z-x, z-y and z-w, while the lower part shows the time series of the system for the Chen system at $a_2=35$, $b_2=3$, $c_2=12$, $d_2=7$, $e=0.5$

Figure 3: Synchronized state $e_{13}$, $e_{22}$, $e_{34}$, $e_{41}$ as shown in (a), (b), (c) and (d) respectively while (e) shows the combined synchronized state of case 1 of non-identical Lorenz (drive) and Chen (response) systems when $a_1=10$, $b_1=8/3$, $c_1=28$, $d_1=2.25$, with initial condition values $x_1=0.1$, $y_1=0.1$, $z_1=0.1$, $w_1=0.1$, the hyperchaotic Chen system (response), parameters selected remain constant at $a_2=35$, $b_2=3$, $c_2=12$, $d_2=7$, $e=0.5$ and variable parameters chosen as $x_2=0.5$, $y_2=$
0.01, \( z_2 = 0.8, w_2 = 0.5 \). The step size was maintained at \( H = 0.005 \) and \( t \geq 10s \).

Figure 4: Combined Synchronized state for cases 2, 3, 4 and 4 non-identical Lorenz (drive) and Chen (response) systems when \( a_1 = 10, b_1 = 8/3, c_1 = 28, d_1 = 2.25 \), with initial condition values \( x_1 = 0.1, y_1 = 0.1, z_1 = 0.1, w_1 = 0.1 \), the hyperchaotic Chen system (response), parameters selected remain constant at \( a_2 = 35, b_2 = 3, c_2 = 12, d_2 = 7, e = 0.5 \) and variable parameters chosen as \( x_2 = 0.5, y_2 = 0.01, z_2 = 0.8, w_2 = 0.5 \). The step size was maintained at \( H = 0.005 \) and \( t \geq 10s \).