Performance of the MAGIC telescopes after the major upgrade

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MAGIC is a system of two Imaging Atmospheric Cherenkov Telescopes located on the Canary island of La Palma, Spain. During summer 2011 and 2012 it underwent a major upgrade. The main subsystems upgraded were the MAGIC-I camera and its trigger system and the readout system of both telescopes. We use observations of the Crab Nebula taken at low and medium zenith angles to assess the key performance parameters of the MAGIC stereo system. For low zenith angle observations, the standard trigger threshold of the MAGIC telescopes is about 50 GeV. The integral sensitivity for point-like sources with Crab Nebula-like spectra above 220 GeV is \((0.66 \pm 0.03)\%\) of Crab Nebula flux in 50 h of observations. The angular resolution, defined as the sigma of a 2-dimensional Gaussian distribution, at energies of a few hundred GeV is below 0.07°, while the energy resolution is around 16%. We investigate the effect of the systematic uncertainty on the data taken with the MAGIC telescopes after the upgrade. We estimate that the systematic uncertainties can be divided in the following components: < 15% in energy scale, 11 − 18% in flux normalization and ±0.15 for the slope of the energy spectrum.

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1. Introduction

MAGIC (Major Atmospheric Gamma Imaging Cherenkov telescopes) is a system of two 17 m diameter Imaging Atmospheric Cherenkov Telescopes (IACTs). The telescopes are located at a height of 2200 m a.s.l. on the Roque de los Muchachos Observatory on the Canary Island of La Palma, Spain (28°N, 18°W). They are used for observations of particle showers produced in the atmosphere by very high energy (VHE, $\gtrsim 30\text{GeV}$) \(\gamma\)-rays. Between summer 2011 and 2012 the telescopes went through a major upgrade (Aleksić et al., 2015a), carried out in two stages. In the first part of the upgrade, in summer 2011, the readout systems of both telescopes were upgraded (Sitarek et al., 2013). The multiplexed FADCs used before in MAGIC-I (Goebel et al., 2007) as well as the Domino Ring Sampler version 2 used in MAGIC-II (DRS2; Tesaro et al., 2009) have been replaced by Domino Ring Sampler version 4 chips (DRS4; Ritt et al., 2010). The switch to DRS4 based readout allowed to eliminate the $\sim 10\%$ dead time present in the previous system due to the DRS2 chip and lowered the electronic noise of the readout. In summer 2012 the second stage of the upgrade followed with an exchange of the camera of the MAGIC-I telescope to a uniformly pixelized one (Aleksić et al., 2015a). Currently each MAGIC camera is equipped with 1039 photomultipliers (PMTs) and has a total field of view of $\sim 3.5^\circ$. The upgrade of the camera allowed to increase the area of the trigger region in MAGIC-I by a factor of 1.7 to the value of 4.8 square degrees.

2. Data sample and analysis

We evaluate the performance of the MAGIC telescopes using data from the Crab Nebula, the so-called standard candle of VHE \(\gamma\)-ray astronomy. The data have been taken between October 2013 and January 2014 in the so-called wobble mode (Fomin et al., 1994), i.e. with the source position offset by $0.4^\circ$ from the center of the camera. This method allows to simultaneously estimate the background from reflected positions in the sky at the same offset. The data are divided into two sub-ranges according to the zenith angle of the observations: $5 - 30^\circ$ (11 h) and $30 - 45^\circ$ (4 h). The data have been analyzed using the standard MAGIC tools: MARS (MAGIC Analysis and Reconstruction Software, Zanin et al., 2013). The analysis chain (extraction of Cherenkov light pulses, calibration to photoelectrons (phe), image cleaning and parametrization, $\gamma$/hadron separation, estimation of energy and arrival direction) is described in detail in Aleksić et al. (2015b).

3. Performance parameters

Here we present the most important performance parameters of the MAGIC telescopes: energy threshold, angular resolution, energy resolution and sensitivity.

3.1 Energy threshold

The energy resolution of an IACT is normally defined as the peak of the differential energy distribution obtained with Monte Carlo (MC) events weighted to a particular spectral shape. In here we assume for the calculation the commonly used spectral index of $-2.6$. 

2
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Figure 1: Left panel: rate of MC $\gamma$-ray events (in arbitrary units) surviving the image cleaning with at least 50 phe for a source with a spectral index of $-2.6$ in the zenith range $< 30^\circ$ (solid line) and $30 - 45^\circ$ (dotted line). Right panel: threshold of the MAGIC telescopes as a function of the zenith angle of the observations at the trigger level (the dotted curve), only events with images that survived image cleaning in each telescope with at least 50 phe (solid line) and with additional cuts of $\text{Hadronness} < 0.5$ and $\theta^2 < 0.03^2$ (dashed line). Both panels reproduced from Aleksić et al. (2015b).

In the left panel of Fig. 1 we show the differential rate plot in two zenith angle ranges for events that survived image reconstruction in both telescopes. For low zenith angle, i.e. $< 30^\circ$, the reconstruction threshold energy is $\sim 70\text{GeV}$. Note however that the peak is broad and extends far to lower energies, allowing to evaluate the performance of the telescopes and to obtain scientific results below such defined threshold. In the right panel of Fig. 1 we show how the energy threshold of the MAGIC telescopes depends on the zenith angle of observations. The threshold is quite stable for low zenith angle observations. It increases rapidly for higher values of zenith angles, due to both the larger absorption of the Cherenkov light in the atmosphere and the dilution of the photons reaching the ground over a larger light pool.

3.2 Energy resolution

We evaluate the performance of the energy reconstruction with $\gamma$-ray MC simulations. The simulations are divided into bins of true energy (5 bins per decade). In each bin we construct a distribution of $(E_{\text{est}} - E_{\text{true}})/E_{\text{true}}$ and fit it with a Gaussian function. The Gaussian function describes well the distribution in the central region, but not at the edges (see Aleksić et al., 2015b). The energy resolution and bias are defined as the standard deviation and the mean value respectively of this fit. Both quantities are plotted as a function of the true energy of the $\gamma$ rays in Fig. 2.

For low zenith angle observations in the energy range of a few hundred GeV the energy resolution is as low as $\sim 15\%$. For higher energies it degrades due to an increasing fraction of truncated images, and showers with high impact parameters as well as worse statistics in the training sample. At low energies the energy resolution is degraded, due to worse precision in the image reconstruction (in particular the impact parameters), and higher internal relative fluctuations of the shower. Above a few hundred GeV the absolute value of the bias is below a few percent. At low energies ($\lesssim 100\text{GeV}$) the estimated energy bias rapidly increases due to the threshold effect. For observations at higher zenith angles the energy resolution is similar. Since an event of the same energy
Figure 2: Energy resolution (solid lines) and bias (dashed lines) obtained from the MC simulations of γ rays at low (red) and medium (blue) zenith angles. Events have been weighted in order to represent a spectrum with a slope of $-2.6$. For comparison, pre-upgrade values from Aleksić et al. (2012) are shown in gray lines. The figure is reproduced from Aleksić et al. (2015b).

observed at higher zenith angle will produce a smaller image, the energy resolution at the lowest energies is slightly worse. On the other hand, at multi-TeV energies, the showers observed at low zenith angle can be partially truncated at the edge of the camera or saturate some of the pixels. Therefore the energy resolution at those energies is slightly better for higher zenith angle observations. On the other hand, as the energy threshold shifts with increasing zenith angle, the energy bias at energies below 100 GeV is much stronger for higher zenith angle observations.

3.3 Angular resolution

Following the approach in Aleksić et al. (2012), we investigate the angular resolution of the MAGIC telescopes using two commonly used methods. In the first approach we define the angular resolution $\Theta_{\text{Gaus}}$ as the standard deviation of a 2-dimensional Gaussian fitted to the distribution of the reconstructed event directions of the γ-ray excess. Such a 2-dimensional Gaussian in the $\theta_x$ and $\theta_y$ space will correspond to an exponential fitting function for $\theta^2$ distribution. The fit is performed in a narrow range, $\theta^2 < 0.025^{\circ2}$. Therefore it is an appropriate quantity for estimating performance of the instrument in case of a search for small extensions (comparable with angular resolution) in VHE γ-ray sources. In the second method we compute an angular distance, $\Theta_{0.68}$, around the source, which contains 68% of the excess events. This method is more sensitive to long tails in the distribution of reconstructed directions. Note that while both numbers assess the angular resolution of the MAGIC telescopes, their absolute values are different. For a purely Gaussian distribution $\Theta_{\text{Gaus}}$ would correspond to only 39% containment radius of γ-rays originating from a point-like source and $\Theta_{0.68} \approx 1.5 \Theta_{\text{Gaus}}$.

We use the low and medium zenith angle samples of the Crab Nebula to investigate the angular resolution. The angular resolution obtained with both methods is shown in Fig. 3. At 250 GeV the angular resolution (computed with the 2D-Gaussian fit) is $0.07^\circ$. It improves with energy, as larger images are better reconstructed, reaching a plateau of $\sim 0.04^\circ$ above a few TeV. The angular resolution improved by about 5-10% after the upgrade. Due to the improvement in angular resolution the small difference between the angular resolution obtained with MC simulations and the Crab
Figure 3: Angular resolution of the MAGIC telescopes after the upgrade as a function of the estimated energy obtained with the Crab Nebula data sample (points) and MC simulations (solid lines). Left panel: 2D Gaussian fit, right panel: 68% containment radius. Red points: low zenith angle sample, blue points: medium zenith angle sample. For comparison the low zenith angle pre-upgrade angular resolution is shown as gray points Aleksić et al. (2012). Reproduced from Aleksić et al. (2015b).

Nebula data, also present in the pre-upgrade data, is slightly more pronounced. The difference of $\sim 10 - 15\%$ is visible at higher energies and corresponds to an additional 0.02° systematic random component (i.e. added in quadrature) between MC and data. Study of the reconstructed source direction spread in each of 10 nightly sub-samples of the Crab Nebula data suggest a similar, $\lesssim 0.02°$, systematic uncertainty on the pointing position (Aleksić et al., 2015b).

3.4 Sensitivity

For a weak source, the significance of an excess of $N_{\text{excess}}$ events over a perfectly-well known background of $N_{\text{bkg}}$ events can be calculated as $N_{\text{excess}}/\sqrt{N_{\text{bkg}}}$. Therefore, one defines the sensitivity $S_{N_{\text{excess}}/\sqrt{N_{\text{bkg}}}}$ as the flux of a source giving $N_{\text{excess}}/\sqrt{N_{\text{bkg}}}=5$ after 50h of effective observation time. For a more realistic estimation of the sensitivity, we apply two additional conditions: $N_{\text{excess}} > 10$ and $N_{\text{excess}} > 0.05N_{\text{bkg}}$. The first condition assures that the Poissonian statistics of the number of events can be approximated by a Gaussian distribution. The second one counteracts the effect of small systematic discrepancies between the ON and OFF distributions, which could mimic a statistically significant signal if the residual background rate is large.

The integral sensitivity of the different phases of the MAGIC experiment for a source with a Crab Nebula-like spectrum are shown in the left panel of Fig. 4. In order to find the optimal cut values in Hadronness ($\gamma$/hadron separation parameter) and $\theta^2$ in an unbiased way, we split the total sample of Crab Nebula data into two independent sub-samples (training and test) of similar size and similar observation conditions. Different energy thresholds are achieved by varying a cut in the total number of photoelectrons of the images (for points < 300GeV) or in the estimated energy of the events (above 300 GeV). For each energy threshold we perform a scan of cuts on the training sub-sample. The pair of cuts which provides the best sensitivity on the training sub-sample is selected and applied to test sample to obtain the final sensitivity value. The threshold itself is estimated as the peak of true energy distribution of MC events with a $-2.6$ spectral slope to which the same cuts were applied.
Figure 4: Integral (left panel) and differential (right panel, 5 bins per decade in energy) sensitivity of the MAGIC telescopes expressed in Crab Units. The sensitivity is defined as the flux of the source in a given energy range for which $N_{\text{excess}}/\sqrt{N_{\text{bkg}}} = 5$ and $N_{\text{excess}} > 10$, $N_{\text{excess}} > 0.05N_{\text{bkg}}$ after 50h of effective time. Left panel: Gray circles: sensitivity of the MAGIC-I single telescope with the Siegen (light gray, long dashed, Albert et al. (2008)) and MUX readouts (dark gray, short dashed, Aleksić et al. (2012)). Black triangles: stereo before the upgrade (Aleksić et al., 2012). Squares: stereo after the upgrade: zenith angle below 30° (red, filled), 30 – 45° (blue, empty). Both panels are reproduced from Aleksić et al. (2015b).

The integral sensitivity evaluated above is valid only for sources with a Crab Nebula-like spectrum. To assess the performance of the MAGIC telescopes for sources with an arbitrary spectral shape, we compute the differential sensitivity, i.e. the sensitivity in narrow bins of energy (5 bins per decade). The differential sensitivity is plotted for low and medium zenith angles in right panel of Fig. 4. The sensitivity computed also according to different definitions is presented in Aleksić et al. (2015b).

The upgrade of the MAGIC-I camera and readout of the MAGIC telescopes has lead to a significant improvement in sensitivity over the whole investigated energy range. The integral sensitivity reaches down to about 0.55% of C.U.\(^1\) around a few hundred GeV in 50 h of observations. The improvement in the performance is especially evident at the lowest energies. In particular, in the energy bin 60-100 GeV, the differential sensitivity decreased from 10.5% C.U. to 6.7% C.U. reducing the needed observation time by a factor of 2.5. Observations at medium zenith angle have naturally higher energy threshold, degrading the performance at the lowest energies. On the other hand such observations provide a small boost in the sensitivity at the TeV energies.

4. Systematic uncertainties

The systematic uncertainties of the IACT technique are a combination of multiple small individual factors which are only known with limited precision, and possibly evolve in time. Most of those factors (e.g. uncertainties connected with the atmosphere, reflectivity of the mirrors, etc.) were not affected by the upgrade. On the other hand the upgrade led to a decrease of the systematic uncertainty connected to background homogeneity (Aleksić et al., 2015b). The total systematic uncertainty can be divided into 3 components: pure flux normalization uncertainty, uncertainty in the energy scale (which can be converted into an additional uncertainty in the flux normalization) and

\(^1\)Crab Unit, i.e. the flux of Crab Nebula, which is a standard candle in VHE γ-ray astronomy.
uncertainty in the spectral index. The final systematic uncertainties have been reevaluated in Aleksić et al. (2015b). The uncertainty on the flux normalization is 18% at low energies ($\lesssim 100\text{GeV}$) and 11% in the energy range of a few hundred GeV. At the highest energies, $\gtrsim 1\text{TeV}$, due to more pronounced MC/data mismatches the systematic uncertainty is a bit higher, namely 16%. The uncertainty of the absolute energy scale of the MAGIC telescopes is below 15%. The systematic uncertainty on the reconstructed spectral slope of the sources is $\pm 0.15$ for observations not strongly dominated by background events.

5. Conclusions

The upgrade of the readout and one of the cameras of the MAGIC telescopes have significantly improved their performance. The trigger threshold for low zenith angle observations is $\sim 50\text{GeV}$. For a Crab-like source, the best performance of the MAGIC telescopes is achieved at energies of a few hundred GeV. At those energies the images are sufficiently large to provide enough information for efficient reconstruction, while the rapidly falling power-law spectrum of the Crab Nebula still provides enough statistics. The energy resolution at these medium energies is as good as 16% with a negligible bias and the angular resolution is $\lesssim 0.07^\circ$. The achieved sensitivity above 220 GeV is $(0.66 \pm 0.03)\%$ of C.U. for 50$h$ of observations. At the lowest energies, below 100 GeV, the performance has improved drastically comparing to the one presented in Aleksić et al. (2012), reducing the needed observation time by a factor of 2.5.

The performance of the MAGIC telescopes at medium zenith angles, $30^\circ - 45^\circ$, is mostly similar to the one at low zenith angles. The higher threshold, however, significantly degraded all the performance parameters below $\sim 200\text{GeV}$. However, for the energies above a few TeV, better performance is achieved with observations at medium zenith angles.

Thanks to the improvement in the performance achieved after the upgrade, the MAGIC telescopes have reached an unprecedented sensitivity, which caused a significant boost in the scientific output of the experiment. In particular the excellent low energy performance allowed to expand the visible VHE $\gamma$-ray sky up to redshift $\sim 0.94$ with the detection of two distant blazars: QSO B0218+357 (Mirzoyan, 2014) and PKS1441+25 (Mirzoyan, 2015).

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