**Chandra** Detection of Intracluster X-Ray Sources in Fornax

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Abstract

Based on archival *Chandra* observations with a total exposure of 1.3 Ms, we study X-ray point sources in the Fornax cluster of galaxies, with the primary aim of searching for intracluster X-ray source populations. We detect 1177 point sources out to a projected radius of ~30 arcmin (~180 kpc) from the cluster center and down to a limiting 0.5–8 keV luminosity of ~3 \times 10^{37} \text{ erg s}^{-1}. We construct the source surface density profile, after excluding sources associated with foreground stars, known globular clusters, ultra-compact dwarfs, and galactic nuclei. From this profile we statistically identify ~183 excess sources that are not associated with the bulk stellar content of the individual member galaxies of Fornax, nor with the cosmic X-ray background. Taking into account Poisson error and cosmic variance, the cumulative significance of this excess is at the $\gtrsim 2\sigma$ level (with a maximum of 3.6$\sigma$) outside three effective radii of the central giant elliptical, NGC 1399. The luminosity function of the excess sources is found to be significantly steeper than that of the GC-hosting sources (presumably low-mass X-ray binaries (LMXBs)), disfavoring the possibility that unidentified GCs are primarily responsible for the excess. We show that a large fraction of the excess can be related to the extended stellar halo of NGC 1399 and/or the diffuse intracluster light, thus providing strong evidence for the presence of intracluster X-ray sources in Fornax, the second unambiguous case for a galaxy cluster after Virgo. Other possible origins of the excess, including supernova-kicked LMXBs and stripped nucleated dwarf galaxies are discussed.

Key words: galaxies: clusters: individual (Fornax) – X-rays: binaries – X-rays: galaxies: clusters

Supporting material: machine-readable table

1. Introduction

X-ray-emitting, close binary systems involving an accreting black hole (BH) or neutron star (NS), were among the first objects discovered in the X-ray sky and are now recognized to be ubiquitous in the local universe, thanks in particular to *Chandra* observations (Weisskopf et al. 2002) of excellent angular resolution and sensitivity. As such, X-ray binaries can serve as a useful tool to study their parent stellar populations (see the review by Fabbiano 2006). It is now a consensus that high-mass X-ray binaries are prevalent in star-forming, typically late-type galaxies, whereas low-mass X-ray binaries (LMXBs) are the predominant population of detected X-ray sources (typically with luminosities $L_X \gtrsim 10^{36} \text{ erg s}^{-1}$) in early-type galaxies (ETGs). This has motivated observational work to calibrate a quasi-linear relation between the total number (or total X-ray luminosity) of LMXBs and host galaxy’s stellar mass (e.g., Gilfanov 2004; Kim & Fabbiano 2004; Zhang et al. 2012).

In recent years, there has been growing evidence that X-ray sources also exist beyond the bulk stellar content of ETGs. Based on *Chandra* observations, Li et al. (2010) identified a significant “excess” of X-ray sources at projected distances of 2–5 optical effective radii of the Sombrero galaxy (M104), which cannot be accounted for by the cosmic X-ray background (CXB) and thus most likely reside in the halo of M104. They postulated that these excess sources, detected at $L_X \approx 10^{37–38} \text{ erg s}^{-1}$, may have a mixed origin. One possibility is that LMXBs dynamically formed in globular clusters (i.e., GC-LMXBs), which tend to have a broader spatial distribution than the field stars. Another possibility is supernova-kicked LMXBs, which have been ejected from the inner galactic regions due to a strong recoil upon the supernova explosion that gives birth to the subsequently accreting NS or BH. Zhang et al. (2013) found a similar excess of halo X-ray sources in a sizable sample of mostly isolated ETGs. The significance of the excess was apparently correlated with both the GC specific frequency and the stellar mass of the host galaxy, supporting an origin of both GC-LMXBs and supernova-kicked LMXBs.

Using *Chandra* observations that sampled 80 intermediate-mass ETGs in the Virgo cluster, Hou et al. (2017) found that an excess of X-ray sources, with $L_X \gtrsim 2 \times 10^{38} \text{ erg s}^{-1}$, exist in the outskirts of these galaxies. On the other hand, they found no significant excess sources of similar luminosities in a control sample of field ETGs. This strongly suggests that at least some of the excess sources are uniquely present in the cluster environment and might be related to the so-called diffuse intracluster light (ICL; e.g., Mihos et al. 2005, 2017). In cluster/group environments, tidal interactions continue to shape and redistribute the stellar content of essentially all member galaxies. In particular, typically low-mass, infalling galaxies are stripped off of their stars, leading to the gradual building of the ICL, a generic process that can begin at the early stage of galaxy clustering (Rudick et al. 2006, 2011). Tidal stripping can also facilitate the growth of an extended stellar halo around the brightest cluster galaxies (BCGs). Being predominantly old stellar populations, the ICL is expected to harbor X-ray sources such as LMXBs. In addition, the stripped relic of nucleated dwarf galaxies, typically of a compact appearance, may also be the host of bright X-ray sources (Hou & Li 2016; Hou et al. 2017). Conversely, a representative
sample of intracluster X-ray sources can serve as a new tool to study the otherwise formidable ICL (Finoguenov et al. 2002), which conventionally requires sensitive optical observations due to its extremely low surface brightness.

The Virgo, however, remains the only cluster in which intracluster X-ray sources have been probed and detected. The Chandra observations utilized by Hou et al. (2017) covered only a small portion (3 deg²) of Virgo, but its large angular size (∼100 deg²) renders a full mapping a challenging, if not infeasible, task with contemporary X-ray telescopes, thus limiting our ability to extend the search for intracluster X-ray sources in Virgo. Located at a distance of 20.0 Mpc (Blakeslee et al. 2009), the Fornax cluster is a dynamically more evolved, more compact and less massive system compared to Virgo. Like Virgo, Fornax has been an important laboratory to explore the physics of hierarchical structure growth. Recent optical surveys including the Hubble Space Telescope ACS Fornax Cluster Survey ( ACSFCS; Jordán et al. 2007) and the Fornax Deep Survey with VST (FDS; Iodice et al. 2016) have significantly advanced our knowledge about the ETGs (Côté et al. 2007; Turner et al. 2012), dwarf galaxies (Venhola et al. 2017), GC populations (Jordán et al. 2015), as well as the ICL (Iodice et al. 2017) of this cluster. In the X-ray band, a survey of the inner ∼30′ region of Fornax has been conducted with Chandra observations (Scharf et al. 2005), which resulted in the detection of more than 700 point-like sources (including CXB sources) against an extended, asymmetric diffuse X-ray emission from the intracluster medium (Ikebe et al. 1996; Jones et al. 1997). The X-ray sources in the BCG, NGC 1399, have been extensively studied with a focus on their connection with GCs and the so-called ultra-compact dwarfs (UCDs; Phillipps et al. 2001). Using archival Chandra observations and assisted with recently advanced knowledge about the ICL in Fornax (see details below), here we search for the expected intracluster X-ray sources. The remainder of this paper is organized as follows. We introduce the Chandra observations and data reduction procedure in Section 2. Our procedures of X-ray source detection and characterization, along with the resultant source catalog, are described in Section 3. In Section 4, we provide strong evidence for the presence of intracluster X-ray sources by exploring the spatial and flux distributions of the detected sources. The possible origins of the intracluster X-ray sources are discussed in Section 5, followed by a summary in Section 6. Quoted errors throughout this work are at the 68% confidence level, unless otherwise noted.

2. Data Preparation

The Fornax cluster, particularly its core region, has been extensively observed by Chandra, chiefly with its Advanced CCD Imaging Spectrometer (ACIS). Initially, we acquired a total of 30 archival Chandra observations, which include 12 ACIS-I and 18 ACIS-S observations. Among these, 10 ACIS-I observations were originally acquired by Scharf et al. (2005), which provide a quasi-uniform exposure of ∼50 ks for the inner ∼30′ of Fornax. Eight additional observations (six ACIS-S and two ACIS-I) were pointed toward the BCG NGC 1399, while another six ACIS-S observations focused on the infalling ETG NGC 1404 (Su et al. 2017), which is centered at a projected radius of ∼9′8 from the center of NGC 1399. These deep exposures enable a good sensitivity for detecting point sources against the strong diffuse X-ray emission from the two galaxies.

We reprocessed the archival level-1 data using CIAO v4.8 and the corresponding calibration files, following the standard procedure.5 We have examined the light curve of each observation and filtered time intervals that suffer from significant particle flares. In particular, ObsID 320 was found to be affected by a strong particle background during its entire 3.5 ks exposure and thus was completely discarded. The background filtering resulted in the final usage of 29 observations with a total cleaned exposure of ∼1300 ks. Across the combined field of view (FoV), the maximum exposure is ∼700 ks (on NGC 1404) and the minimum is ∼40 ks. A log of the 29 observations is given in Table 1.

Following Hou et al. (2017), we produced count maps and exposure maps, at the natal pixel scale of 0.′492. for each ObsID in three energy bands: 0.5–2 (S-band), 2–8 (H-band), and 0.5–8 (F-band) keV. The exposure maps were weighted by an assumed incident spectrum of an absorbed power law, with photon index of 1.7, and absorption column density $N_H = 10^{21}$ cm$^{-2}$. This latter value is higher than the Galactic foreground absorption column (∼1.5 × 10$^{20}$ cm$^{-2}$), but allows for some intrinsic absorption in LMXBs (e.g., Luan et al. 2018). To ensure an optimal sensitivity for source detection, we included only data from CCDs I0, I1, I2, and I3 for the ACIS-I observations, and CCDs S2 and S3 for the ACIS-S observations. We also generated for each band and each ObsID maps of the point-spread function (PSF), at a given enclosed count fraction (ECF), using the same spectral weighting as for the exposure map. Lastly, we reprojected the individual count maps or exposure maps to a common tangential point, i.e., the nucleus of NGC 1399 ([R.A., decl.] = [54.620941, −35.450657]), to produce a combined image. The PSF maps were similarly combined, with weights according to the local effective exposure.

Figure 1(a) presents the 0.5–8 keV flux image of the combined FoV. A close-up view of the vicinity of NGC 1399 and NGC 1404, where both strong diffuse X-ray emission and numerous discrete sources are clearly present, is displayed in Figures 1(b) and (c), respectively.

3. X-Ray Source Detection

Our source detection and source characterization procedures are as detailed in Zhu et al. (2018). Here we briefly outline the key steps and results.

(i) We performed source detection for each of the three energy bands over the combined FoV, using the CIAO tool wavdetect. The combined exposure map and 50%-ECF PSF map were supplied to the detection process. A false-positive probability threshold of 10$^{-6}$ was adopted. This resulted in a raw list of 1008 sources in S-band, 758 sources in H-band, and 1248 sources in F-band. At this point, we also produced a sensitivity map, which records the detection limit at each pixel, following the recipe of Kashyap et al. (2010) and according to the chosen false-positive threshold.

5 http://cxc.harvard.edu/ciao/
Table 1
Log of Chandra Observations

| ObsID | Instrument | R.A. | Decl. | Exposure | Start Date |
|-------|------------|------|-------|----------|------------|
| 624   | ACIS-S     | 03 39 34.70 | −35 25 50.00 | 43.6 | 1999 Dec 15 |
| 319   | ACIS-S     | 03 38 29.40 | −35 27 00.40 | 56.0 | 2000 Jan 18 |
| 239   | ACIS-I     | 03 38 29.40 | −35 27 00.40 | 3.6 | 2000 Jan 19 |
| 2942  | ACIS-S     | 03 38 52.00 | −35 35 34.00 | 29.2 | 2003 Feb 13 |
| 4168  | ACIS-I     | 03 36 59.85 | −35 29 39.84 | 45.6 | 2003 May 20 |
| 4169  | ACIS-I     | 03 37 09.50 | −35 44 00.06 | 37.4 | 2003 May 21 |
| 4171  | ACIS-I     | 03 38 11.20 | −35 41 54.20 | 45.0 | 2003 May 23 |
| 4170  | ACIS-I     | 03 37 14.40 | −35 13 27.84 | 40.6 | 2003 May 24 |
| 4172  | ACIS-I     | 03 38 25.56 | −35 25 42.60 | 44.5 | 2003 May 26 |
| 4173  | ACIS-I     | 03 38 37.01 | −35 09 27.72 | 45.1 | 2003 May 26 |
| 4174  | ACIS-I     | 03 38 49.58 | −35 34 36.34 | 45.7 | 2003 May 28 |
| 4175  | ACIS-I     | 03 39 30.51 | −35 45 22.03 | 54.6 | 2003 May 29 |
| 4176  | ACIS-I     | 03 39 44.60 | −35 29 08.92 | 46.0 | 2003 May 31 |
| 4177  | ACIS-I     | 03 39 55.80 | −35 12 56.16 | 38.9 | 2003 Jun 1 |
| 3949  | ACIS-S     | 03 40 10.50 | −35 37 38.00 | 54.6 | 2003 Oct 21 |
| 9798  | ACIS-S     | 03 38 51.00 | −35 34 31.00 | 18.3 | 2007 Dec 24 |
| 9799  | ACIS-S     | 03 38 51.00 | −35 34 31.00 | 21.3 | 2007 Dec 27 |
| 9530  | ACIS-S     | 03 38 29.00 | −35 27 01.40 | 59.4 | 2008 Jun 8 |
| 14527 | ACIS-S     | 03 38 29.10 | −35 27 03.00 | 27.8 | 2013 Jul 1 |
| 16639 | ACIS-S     | 03 38 29.10 | −35 27 03.00 | 29.7 | 2014 Oct 12 |
| 16231 | ACIS-S     | 03 39 01.20 | −35 35 18.60 | 60.5 | 2014 Oct 20 |
| 17541 | ACIS-S     | 03 39 01.20 | −35 35 18.60 | 24.7 | 2014 Oct 23 |
| 16334 | ACIS-S     | 03 39 22.32 | −35 38 42.00 | 90.9 | 2014 Oct 30 |
| 17540 | ACIS-S     | 03 39 01.20 | −35 35 18.60 | 28.5 | 2014 Nov 2 |
| 16233 | ACIS-S     | 03 39 04.08 | −35 35 31.20 | 98.8 | 2014 Nov 9 |
| 17548 | ACIS-S     | 03 39 04.08 | −35 35 31.20 | 48.2 | 2014 Nov 11 |
| 16232 | ACIS-S     | 03 39 04.08 | −35 35 31.20 | 69.1 | 2014 Nov 12 |
| 17549 | ACIS-S     | 03 38 51.26 | −35 38 19.28 | 61.7 | 2015 Mar 28 |
| 14529 | ACIS-S     | 03 38 29.10 | −35 27 03.00 | 31.6 | 2015 Nov 6 |

Note. (1) Observation ID. (2) Instrument. (3)–(4) Epoch 2000 Coordinates of the aim-point. (5) Cleaned exposure, in units of ks. (6) Start date of the observation, in the form of yyyy-mm-dd.

(ii) We refined the source centroids output from wavdetect, using a maximum likelihood method (Boese & Doebereiner 2001) that iterates over the recorded positions of the individual counts within the 90% enclosed counts radius (ECR) as default. A 50% ECR was adopted for a small fraction of sources to avoid confusion with closely neighboring sources. The position uncertainty (PU) at the 68% confidence level was estimated following the empirical relation between PU, source counts, and source position in terms of the off-axis angle (Kim et al. 2007).

(iii) We then performed source photometry to derive the exposure- and PSF-corrected photon flux. A circular aperture was chosen, with a default 90% ECR to extract the source counts; for those sources subject to crowding, the 50% ECR was again adopted. The background extraction regions were typically 2–4 times the 90% ECR, excluding pixels falling within 2 times the 90% ECR of any neighboring sources. A Bayesian approach was employed to calculate the photon fluxes and bounds, which takes care of the Poisson statistics at the low-count regime (Park et al. 2006). We also calculated the hardness ratio, defined as $HR = (S_{0.5-8} - S_{0.5-2})/(S_{2-8} + S_{0.5-2})$, where $S_{0.5-2}$ and $S_{0.5-8}$ are the photon flux of the $S$-band and $H$-band, respectively.

(iv) We calculated the binomial no-source probability ($P_B$; Weisskopf et al. 2007) to filter spurious sources due to background fluctuation. Any source with $P_B > 0.01$ is considered spurious and excluded from further analysis, while the remaining are considered genuine point sources.

(v) A cross-matching method (Hong et al. 2009) was then employed to identify the same source detected in more than one band. The relative distance ($d_r$) of two sources from two different bands, is defined as the ratio of the angular offset between the source centroids to the quadratic sum of the respective PU. We required $d_r < 3.0$ for a true source pair. With this criterion, we estimate 8% random matches out of the 881 (659) $S/F$ ($H/F$) pairs, i.e., $\sim 1\%$ false matches.

At this point, we arrive at a list of 1279 independent sources, among which 1177 are detected in the $F$-band, 926 in the $S$-band, and 716 in the $H$-band. For future reference, we present in Table 2 a catalog of basic source properties, including centroid position, observed photon flux in the three bands, 0.5–8 keV unabsorbed energy flux, and the hardness ratio. To derive the 0.5–8 keV unabsorbed flux, we have adopted a photon flux-to-energy flux conversion factor of $3.64 \times 10^{-9}$ erg ph$^{-1}$ according to the assumed incident source spectrum (Section 2).

In Figure 2, we plot the observed 0.5–8 keV photon flux ($S_{0.5-8}$) versus the projected distance ($R$) from the cluster center (here defined as the center of NGC 1399) for the $F$-band sources. For comparison, we also show the azimuthally averaged detection limit as a function of $R$, which is derived...
Figure 1. (a) The mosaic 0.5–8 keV flux image of the Fornax cluster, smoothed by a 2 pixel Gaussian kernel. The center of NGC 1399 and NGC 1404 is marked by a red “X” and a blue “+,” respectively. The locations of other FCC galaxies within the FoV (Table 3) are marked by black crosses. A close-up view of the vicinity of NGC 1399 and NGC 1404 is shown in panels (b) and (c), respectively, where deep exposures are available. The small black circles highlight X-ray sources associated with known globular clusters. See Section 3 for details.
from the sensitivity map. A limiting photon flux of \( \sim 3 \times 10^{-7} \text{ ph cm}^{-2} \text{ s}^{-1} \) is reached within \( R \approx 4' \), where a deep exposure competes with the strong diffuse emission from NGC 1399. It is noteworthy that the global limiting flux, \( \sim 1.5 \times 10^{-7} \text{ ph cm}^{-2} \text{ s}^{-1} \), is actually achieved around NGC 1404. Outside NGC 1399 and NGC 1404, and out to \( R \approx 25' \), the median detection limit is roughly leveled at \( (1-2) \times 10^{-6} \text{ ph cm}^{-2} \text{ s}^{-1} \) due to the highly uniform effective exposure. Figure 3 shows HR versus \( S_{0.5-8} \) for the F-band sources, along with predicted hardness ratios of certain absorbed power laws. It can be seen that the majority of sources exhibit HR values consistent with them being LMXBs and/or background active galactic nuclei (AGNs), i.e., with photon-indices of 1–2. A small number of hard sources (HR \( \gtrsim 0.5 \)) might be heavily obscured AGNs.

### 3.1. Optical Counterpart

Our last step before turning to search for intracluster X-ray sources, involves the identification of a priori irrelevant sources on a best-effort basis. First, we identify possible Galactic foreground X-ray sources by cross-correlating with the USNO-B catalog (Monet et al. 2003) for optical sources with significant proper motions. This results in 18 pairs for a matching radius of 1'. Second, we identify any source that is located within 3\( \arcmin \) from the optical nucleus of any member galaxy of Fornax. Within the Chandra FoV, there are 29 known member galaxies according to the Fornax Cluster Catalog (FCC; Ferguson 1989), whose positions are marked in Figure 1(a). Among them, a nuclear X-ray source is found in NGC 1381 and NGC 1387; in each case of NGC 1399 and NGC 1404, two sources satisfy our matching criterion, and the one with the smaller offset is chosen as the nuclear source. The basic information of the 29 FCC galaxies is summarized in Table 3.

In addition, we identify X-ray sources positionally coincident with known GCs in Fornax. First, we consult with the ACSFCS GC catalog (Jordán et al. 2015), which, with a limiting g-band magnitude of 26.3, is expected to contain \( >90\% \) of the GC population that falls within the ACS fields. We select GC candidates with \( p_{GC} \geq 0.5 \) in the ACSFCS catalog and adopt a matching radius of 1'' between the X-ray and optical centroids, which results in 134 pairs. Since the ACSFCS fields only cover a small fraction of the Chandra FoV, we also incorporate the GC catalog based on the FDS (Cantiello et al. 2018), which fully overlaps with the Chandra FoV and reaches a depth of 14.5. A matching radius of 1'' results in 74 pairs from the FDS catalog. We then cross-correlate two additional GC catalogs, which are intermediate between the ACSFCS and FDS catalogs in terms of sky coverage and sensitivity. A matching radius of 1'' results in 118 pairs from the catalog of Kim et al. (2013), which is

### Table 2: Catalog of Detected X-Ray Sources in the Fornax Cluster

| No. | R.A.  | Decl. | Pos. Err | \( S_{0.5-8} \) | \( S_{2-8} \) | \( S_{0.5-8} \) | \( F_{0.5-8} \) | HR | Note |
|-----|------|-------|----------|----------------|------------|-------------|-------------|-----|------|
| 1   | 54.07722 | −35.48106 | 1.60     | <5.9           | <22.0      | 14.8         | 5.0         | 53.9 | 0.42 |
| 2   | 54.08801 | −35.55983 | 0.42     | 216.3        | 54.3       | 269.3       | 1.8         | 980.3 | 0.06 |
| 3   | 54.11255 | −35.51206 | 0.75     | 14.3         | <16.2      | 23.0         | 0.5         | 83.6 | 0.20 |
| 4   | 54.12001 | −35.37550 | 1.06     | 35.4         | <31.4      | 54.1         | 2.9         | 197.0 | 0.15 |
| 5   | 54.12192 | −35.54933 | 0.60     | 39.2         | <14.4      | 46.3         | 1.8         | 167.8 | 0.14 |
| 6   | 54.13016 | −35.44953 | 0.79     | 10.7         | <15.4      | 18.8         | 0.5         | 68.4 | 0.28 |
| 7   | 54.13259 | −35.29532 | 1.05     | 167.6        | <42.7      | 143.1        | 1.8         | 521.0 | 0.12 |
| 8   | 54.13600 | −35.29274 | 1.08     | <114.8       | <59.0      | 114.7        | 3.0         | 417.7 | 0.14 |
| 9   | 54.13765 | −35.59933 | 0.85     | 43.2         | <25.7      | 56.8         | 0.6         | 206.8 | 0.12 |
| 10  | 54.14037 | −35.41325 | 0.56     | 35.1         | 22.3        | 57.2         | 0.8         | 208.4 | 0.16 |

Note. (1) Source ID, in order of increasing R.A. (2)–(3) R.A. and decl. (J2000) of source centroid. (4) Position uncertainty in arcseconds. (5)–(7) The 0.5–2 keV, 2–8 keV, and 0.5–8 keV photon flux, in units of \( 10^{-7} \text{ ph cm}^{-2} \text{ s}^{-1} \). (8) The 0.5–8 keV energy flux, in units of \( 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1} \), converted from 0.5–8 keV photon flux by assuming an absorbed power-law spectrum with photon index of 1.7 and column density \( 10^{21} \text{ cm}^{-2} \). (9) Hardness ratio between the 0.5–2 keV and 2–8 keV bands. (10) “F” denotes foreground stars, “G” denotes sources spatially coincident with known GCs, “N” denotes nuclear source of a member galaxy, and “U” denotes sources spatially coincident with known UCDs. Quoted errors are at the 1\( \sigma \) confidence level, while 3\( \sigma \) upper limits are given in the case of nondetection in a given band.

(This table is available in its entirety in machine-readable form.)

![Figure 2. Observed 0.5–8 keV photon flux vs. the projected distance from the center of NGC 1399, for sources detected in the F-band. The red solid and blue dashed curves denote the azimuthally minimum and median detection sensitivity, respectively. The “valley” between 6‘ and 13‘ in the minimum sensitivity is due to the deep exposure toward the vicinity of NGC 1404.](image-url)
Figure 3. Hardness ratio vs. observed 0.5–8 keV photon flux, for sources detected in the F-band. The y-axis on the right shows the corresponding unabsorbed luminosity assuming the distance of Fornax. The vertical dashed lines represent predicted hardness ratios by various absorbed power-law spectral models. The cyan dotted–dashed line represents column density $N_H = 1.5 \times 10^{20}$ cm$^{-2}$ and photon index of 1.7; the red, blue, and green dashed lines represent $N_H = 10^{21}$ cm$^{-2}$ and photon index of 2.0, 1.7, and 1.4, respectively; the magenta dotted line represents $N_H = 10^{22}$ cm$^{-2}$ and photon index of 1.7. The three error bars, from top to bottom, indicate the median errors of sources with $S_{0.5-8} > 10^{-5}$ ph cm$^{-2}$ s$^{-1}$, $10^{-6}$ ph cm$^{-2}$ s$^{-1} < S_{0.5-8} < 10^{-5}$ ph cm$^{-2}$ s$^{-1}$, and $S_{0.5-8} < 10^{-6}$ ph cm$^{-2}$ s$^{-1}$.

Based on Blanco-4 m observations that cover a 36′ × 36′ field around NGC 1399 down to a limiting $U$-band magnitude of 24.4; the same matching radius finds 130 pairs from the catalog of Paolillo et al. (2011), which is based on ACS V-band mapping of a ∼10′ × 10′ field centered at NGC 1399.

In total, we have identified 270 independent X-ray sources associated with Fornax GCs, which are presumably GC-LMXBs. By artificially shifting the X-ray centroids by 10″ in both directions of R.A. and decl., we estimate that overall only a few percent of all GC-LMXB pairs can be random matches, although the amount increases to ∼20% in the inner few arcminutes of NGC 1399, where the surface density of both X-ray and optical sources is high. On the other hand, due to the heterogeneous nature of the GC catalogs in use, it is not straightforward to assess the completeness of GC-LMXBs. Empirically, the vast majority of GC-LMXBs are found in massive GCs (e.g., Li et al. 2010; Hou & Li 2016). Indeed, among the 134 ACSFCS GC-LMXB pairs, ∼73% are found in GCs brighter than the so-called turnover magnitude (∼24.0 mag of g-band at the distance of Fornax; Villegas et al. 2010). We can have a rough estimate of potentially missing GC-LMXBs in the other catalogs, by scaling the number of pairs found at a certain magnitude range with respect to the ACSFCS catalog. In particular, the FDS catalog, which covers the entire Chandra FoV and just reaches the turnover magnitude, is expected to miss $\lesssim$20 fainter, low-mass GCs that might have an X-ray counterpart.

We further cross-correlate the detected X-ray sources with catalogs of Fornax UCDs (Gregg et al. 2009; Voggel et al. 2016), finding 15 pairs with a matching radius of 1″. It turns out that among them, 13 have already been classified as a GC. This is not totally surprisingly, because some UCDs can appear as “giant” GCs in optical images. The other two UCDs (source ID 326 and 717 in Table 2) are new identifications.

The above X-ray sources with an identified optical counterpart is denoted in the last column of Table 2 (“F” for foreground sources, “N” for nuclear sources, “G” for GC-LMXBs, and “U” for UCDs).

4. Analysis: Probing Intracluster X-Ray Sources

4.1. Spatial Distribution

The presence of intracluster X-ray sources, defined as sources that are located in Fornax but outside the main stellar content of the member galaxies, is best probed by the source spatial distribution (Hou et al. 2017). We focus on the sources detected in the 0.5–8 keV band, having excluded foreground sources, nuclear sources, UCDs, and probable GC-LXMBs as identified in the previous section. It is noteworthy that some GC-LMXBs can belong to the putative intracluster stellar populations. Indeed, in the case of Virgo, ∼30% of the “excess” X-ray sources can be attributed to GC-LMXBs (Hou et al. 2017). Here we preclude known GC-LMXBs (Section 3.1) to focus on the “field” population. The relevance of GC-LMXBs to the intracluster populations will be further addressed in Section 5.

In addition, we preclude sources that are most likely associated with the member galaxies other than NGC 1399. Such sources are defined as those falling within three times the optical effective radius ($R_e$) of a given host galaxy (Table 3), a
Table 3

| Galaxy Name | Other Name | R.A. (1) | Decl. (2) | R_c (3) | Source ID |
|-------------|------------|----------|-----------|---------|-----------|
| FCC 170     | NGC 1381   | 54.1317  | -35.2953 | 12.9    | 7         |
| FCC 171     |            | 54.1557  | -35.3846 | 18.0    |           |
| FCC 175     |            | 54.1779  | -35.4341 | 7.7     |           |
| FCC 182     |            | 54.2261  | -35.3729 | 11.4    |           |
| FCC 184     | NGC 1387   | 54.2369  | -35.5066 | 50.1    | 63        |
| FCC 188     |            | 54.2689  | -35.5886 | 12.1    |           |
| FCC 190     | NGC 1380B  | 54.2877  | -35.1937 | 16.3    |           |
| FCC 191     |            | 54.2913  | -35.3854 | 6.3     |           |
| FCC 193     | NGC 1389   | 54.2987  | -35.7443 | 20.1    |           |
| FCC 194     |            | 54.3243  | -35.6972 | 6.4     |           |
| FCC 196     |            | 54.3909  | -35.8277 | 10.2    |           |
| FCC 197     |            | 54.4203  | -35.2948 | 6.3     |           |
| FCC 202     | NGC 1396   | 54.5267  | -35.4383 | 9.8     |           |
| FCC 207     |            | 54.5797  | -35.1274 | 8.5     |           |
| FCC 208     |            | 54.5778  | -35.5209 | 11.7    |           |
| FCC 211     |            | 54.5889  | -35.2582 | 5.6     |           |
| FCC 213     | NGC 1399   | 54.6215  | -35.4506 | 49.1    | 527       |
| FCC 214     |            | 54.652   | -35.8324 | 6.2     |           |
| FCC 215     |            | 54.6562  | -35.7555 | 7.4     |           |
| FCC 218     |            | 54.6887  | -35.2645 | 7.6     |           |
| FCC 219     | NGC 1404   | 54.7171  | -35.5938 | 20.0    | 841       |
| FCC 220     |            | 54.7292  | -35.2352 | 5.6     |           |
| FCC 222     |            | 54.8053  | -35.3697 | 14.5    |           |
| FCC 223     |            | 54.8311  | -35.7234 | 16.6    |           |
| FCC 227     |            | 54.9587  | -35.5211 | 7.3     |           |
| FCC 228     |            | 54.9636  | -35.3206 | 9.2     |           |
| FCC 229     |            | 54.9799  | -35.6609 | 6.8     |           |
| FCC 235     | NGC 1427A  | 55.0383  | -35.6231 | 36.3    |           |
| FCC 241     |            | 55.097   | -35.2744 | 15.9    |           |

Note. (1) Member galaxies in the Fornax Cluster Catalog (Ferguson 1989). (2) Other names of the galaxy. (3)-(4) J2000 coordinates of the galactic center. (5) Effective radius in arcseconds; the value for NGC 1399 is adopted from Iodice et al. (2016). (6) ID of the identified nuclear X-ray source, as listed in Table 2.

Figure 4(a). It is obvious that the observed and predicted source profiles are highly consistent with each other at $R \gtrsim 20'$. The central component, presumably field-LMXBs, is expected to follow the starlight distribution of NGC 1399, which can be characterized by a Sérsic law,

$$
\mu(R) = \mu_e + k_e \left( \frac{R}{R_e} \right)^{1/n} - 1,
$$

where the parameters were determined from the FDS (Iodice et al. 2016): $\mu_e = 21.5$ mag arcsec$^{-2}$, $R_e = 49.1$ arcsec, and $n = 4.5$ ($k_e = 2.17 n - 0.355$; Caon et al. 1993) in the g-band.
We note that these parameters were derived by excluding the stellar core ($R < 0'1$) of NGC 1399, hence we also neglect the few X-ray sources at $R < 0'1$ in the following comparison. To convert the starlight into stellar mass, we adopt the color-dependent mass-to-light ratio of Bell et al. (2003), finding $M/L \approx 5.2$ for g-band solar absolute magnitude of 5.15 and a quasi-uniform color $g - i = 1.2$ (Iodice et al. 2016). The Sérsic profile is then normalized to match the observed X-ray source profile within $R \lesssim 3R_e \approx 2'5$ (blue dotted curve in Figure 4(a)). We have applied the field-LMXB luminosity function (LF) of Zhang et al. (2011) to correct the Sérsic profile for the position-dependent detection incompleteness. The normalization ($\epsilon_X$) of the Sérsic profile is found to be $\approx 4.3$ sources per $10^{10} M_\odot$ (above a reference luminosity $L_{\text{ref}} = 5 \times 10^{37}$ erg s$^{-1}$). This value is in good agreement with the range of specific number of LMXBs above a similar limiting luminosity found in ETGs (2.8–5.4 sources per $10^{10} M_\odot$, Zhang et al. 2011, 2012).

The combination of the above two components, however, cannot fully account for the observed source profile; an excess is clearly seen over $3' \lesssim R \lesssim 15'$ (Figure 4(b)). We stress that adjusting $\epsilon_X$, the only free parameter, does not eliminate the excess, due to the steeply declining Sérsic profile. We define the excess as $N_{\text{excess}} = N_{\text{obs}} - N_{\text{ser}} - N_{\text{CXB}}$, and quantify its cumulative significance as:

$$
\text{SIG} = \frac{N_{\text{obs}} - N_{\text{ser}} - N_{\text{CXB}}}{\sqrt{N_{\text{obs}}^2 \frac{GM_\odot^2}{L_{\text{ref}}^2} + N_{\text{CXB}}^2 \frac{GM_\odot^2}{L_{\text{ref}}^2}}},
$$

where $N_{\text{obs}}$ is the number of observed sources, and $N_{\text{ser}}$ and $N_{\text{CXB}}$ are the number of predicted sources in the Sérsic and CXB components, respectively. We take into account the Poisson variance $\sigma_p^2 = 1/N_{\text{obs}}$ and the cosmic variance $\sigma_r^2$, but neglect the uncertainty in $N_{\text{ser}}$ which is small compared to the other two errors over the radial range where the excess is concerned. The cosmic variance can be estimated as (Lahav & Saslaw 1992),

$$
\sigma_c^2 = \frac{1}{\Omega^2} \int_0^\infty \omega(\theta) d\Omega_1 d\Omega_2 = C_\gamma \theta_0^{-1} \Theta^{-\gamma},
$$

where $\omega(\theta) = (\theta/\theta_0)^{1-\gamma}$ is a power-law angular correlation function (Peebles 1980), $\theta_0 \approx 0'00214$ is the correlation length (Ebrero et al. 2009), $\Omega = \Theta^2 \approx 0.66$ deg$^2$ is the size of the FoV, and the numerical factor $C_\gamma \approx 2.25$ for the canonical value of $\gamma = 1.8$ (Peebles 1980). This results in $\sigma_c \approx 0.14$. As shown in Figure 4(e), the excess has a cumulative significance $\geq 2\sigma$ everywhere beyond $R = 3 R_e$, reaching a maximum of $3.6\sigma$ at $7'5 \lesssim R \lesssim 12'5$. 
Based on the FDS, Iodice et al. (2016) detected an extended and diffuse stellar halo around NGC 1399, which is distinct from the Sérsic component and can be traced out to a projected distance of ~190 kpc. The azimuthally averaged surface brightness profile of this halo was described by an exponential law:

\[ \mu(R) = \mu_0 + 1.086(R/R_0), \]  

where in the g-band \( \mu_0 = 23.4 \) mag arcsec\(^{-2} \) and \( R_0 = 292 \) arcsec. It is conceivable that the at least part of the excess X-ray sources can be attributed to this stellar halo. We test this possibility by adding the above exponential component to our spatial model, again correcting for detection incompleteness with respect to the empirical LF of field-LMXBs and adopting the same \( M/L \) as for the Sérsic component. We then fit the observed source profile using the Sérsic+exponential+CXB model over the range of 0'/1 < R < 27'/5, the outer boundary chosen as the radius beyond which the azimuthal coverage of the Chandra FoV drops below 50%. Initially, we assume that the same normalization applies to both the Sérsic and exponential components, finding \( \epsilon_X = 4.3 \pm 0.4 \) sources per \( 10^{10} M_\odot \) (above \( L_{\text{ref}} \)) and \( \chi^2/\text{d.o.f.} = 64.2/54 \). The fitted exponential component, with a total number of \( N_{\text{exp}} = 109 \) sources, is plotted as the red dotted curve in Figure 4(a), while the combined model is plotted as the solid curve. This simple exercise suggests that the exponential halo can be responsible for a substantial fraction of the excess sources. On the other hand, as indicated in Figures 4(c) and (d), even after accounting for the exponential component, there is still a significant residual (i.e., \( N_{\text{obs}} - N_{\text{star}} - N_{\text{exp}} - N_{\text{CXB}} \)) of ~74 sources (or 40% of all excess sources). We also test the case that the two spatial components have different normalizations, obtaining \( \epsilon_X \) (Sérsic) = 5.3 \pm 1.0 sources per \( 10^{10} M_\odot \) and \( \epsilon_X \) (exponential) = 3.5 \pm 0.9 sources per \( 10^{10} M_\odot \), with \( \chi^2/\text{d.o.f.} = 62.9/53 \). This indicates no significant difference in the source abundance between the two components. The role of the stellar halo of NGC 1399 in the excess sources will be further addressed in Section 5.

### 4.2. Luminosity Function

The source properties can be further constrained by examining their LF. In the case of Virgo, however, the moderate number and limited flux range of the excess sources did not permit a meaningful LF analysis (Hou et al. 2017). Here we consider three sets of F-band sources: (i) “Fornax-field”
Figure 5. Observed luminosity function of F-band sources, adaptively binned to have a minimum of nine sources per flux bin. Black (Fornax-field): field sources (i.e., excluding GC-LMXBs and sources associated with member galaxies) detected at $4' < R < 15'$. Blue (Fornax-GC): GC-LMXBs detected at $4' < R < 15'$. Orange (NGC 1399-field): field sources detected at $R < 3 R_e$. The fitted model, modified by detection incompleteness and Eddington bias, is plotted as the solid curve. For Fornax-field, the model includes a power law (dashed curve) and a fixed CXB component (dotted curve). For NGC 1399-field and Fornax-GC, only the power law is adopted.

![Observed Luminosity Function](image)

refers to sources detected between $4' < R < 15'$, excluding known GC-LMXBs and UCDs but necessarily including CXB contribution; (ii) “NGC 1399-field” refers to field sources detected at $R < 3 R_e$, where CXB contribution is negligible (Figure 4), and (iii) “Fornax-GC” refers to GC-LMXBs detected between $4' < R < 15'$, but excluding four objects that are also classified as UCDs.6

The three observed LFs are shown in Figure 5, which consists of 302 (Fornax-field), 85 (NGC 1399-field), and 82 (Fornax-GC) sources, respectively. We fit the LFs using a canonical power-law model, $dN/dL \propto L^{-\alpha}$. This model is corrected for incompleteness and Eddington bias, following the procedure of Wang (2004) and Zhu et al. (2018) and taking into account the source spatial distribution. To zeroth order, a flat distribution is assumed for Fornax-field and Fornax-GC, and a Sérsic profile is assumed for NGC 1399-field (Section 4.1). For Fornax-field, we also include a fixed component, according to the log $N$–log $S$ relation of Georgakakis et al. (2008), to account for the CXB contribution (dotted curve in Figure 5). The Fornax-field LF shows a clear excess above the CXB level at photon fluxes $\lesssim 2 \times 10^{-6}$ ph cm$^{-2}$ s$^{-1}$. The best-fit power-law slope is obtained by minimizing the C-statistic (Cash 1979). As summarized in Table 4, the fit result indicates a rather steep LF of the Fornax-field ($\alpha = 2.54^{+0.38}_{-0.29}$, 90% confidence level), which is statistically consistent with that of the NGC 1399-field ($\alpha = 2.35^{+0.22}_{-0.20}$). On the other hand, the Fornax-GC LF is significantly flatter, with $\alpha = 1.94^{+0.15}_{-0.14}$. That the GC-LMXBs have a flatter LF than their field counterparts is consistent with previous work (e.g., Kim et al. 2009; Zhang et al. 2011).

We note that in the above analysis we have ignored $\sim$100 sources detected only in the $S_4$- or $H$-band, chiefly to minimize the uncertainty in the CXB contribution and the LF. Some of these sources can also belong to the intracluster populations.

5 According to Hou & Li (2016), the X-ray emission from UCDs is most likely due to dynamically formed LMXBs, as in the case of GCs. Here we conservatively distinguish the X-ray counterparts of UCDs and GCs.

6 In Hou et al. (2017), the excess was identified without excluding GC-LMXBs known a priori. The current choice of precluding GC-LMXBs to our best knowledge therefore means a more stringent definition of excess sources.

5. Discussion

We have investigated the spatial and flux distributions of X-ray point sources in Fornax, out to a projected distance of $\sim 30'$ ($\sim 180$ kpc) from the cluster center. After accounting for sources of well-established origins, i.e., those associated with the member galaxies (in particular the BCG), known GCs and the CXB, significant excess is found over a substantial radial range (Figure 4). Below we discuss the possible nature of these excess sources, in close comparison with their counterparts previously identified in the Virgo cluster (Hou et al. 2017).

The first and obvious possibility is that some of the excess sources are still GC-LMXBs, despite our comprehensive effort in identifying them. Indeed, in the case of Virgo, among the $\sim 120$ excess sources with $L_X \gtrsim 10^{38}$ erg s$^{-1}$, $\sim 30\%$ can be attributed to GC-LMXBs7 (Hou et al. 2017). In Fornax, we found GC-LMXBs out to a radius of 25'/6, and GCs have been detected even beyond the edge of our FoV (Bassino et al. 2006; Cantiello et al. 2018). As addressed in Section 3.1, the majority of the GC-LMXBs are found in bright, massive GCs, and only $\lesssim 20$ X-ray sources associated with faint, undetected GCs are likely to have escaped our identification. Moreover, the LF of Fornax-GC sources is much flatter than that of the Fornax-field sources (Section 4.2), suggesting that the two sets of sources have different origins. Hence we conclude that any unidentified GC-LMXBs would have only a minor contribution to the excess sources.

The second plausible scenario to consider is supernova-kicked LMXBs, favorably consisting of an accreting NS. If the NS received a kick velocity greater than the host galaxy’s escape velocity (on the order of 100 km s$^{-1}$), which is due to its parent supernova explosion, and the binary system in which the NS resides survived this kick, an NS-LMXB might be later found at the galaxy outskirts or the intracluster space (Brandt & Podsiałowski 1995; Zuo et al. 2008). While such a scenario could be generic, supernova-kicked LMXBs were considered to have a negligible contribution to the excess sources found in Virgo (Hou et al. 2017), because the latter were detected with $L_X \gtrsim 2 \times 10^{38}$ erg s$^{-1}$, i.e., above the Eddington limit for NS-LMXBs. Given the current detection limit of $\sim 3 \times 10^{37}$ erg s$^{-1}$ (Figure 3), supernova-kicked LMXBs are likely present in the excess sources. Zhang et al. (2013) proposed that

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Table 4

| Source           | $\alpha$       | $N(>L_{\text{ref}})$ | C/d.o.f. |
|------------------|----------------|----------------------|----------|
| Fornax-field     | 2.54$^{+0.38}_{-0.29}$ | 306                  | 41.7/23  |
| Fornax-GC        | 1.94$^{+0.15}_{-0.14}$ | 199                  | 5.3/7    |
| NGC 1399-field   | 2.35$^{+0.22}_{-0.20}$ | 126                  | 11.8/6   |

Note. The luminosity functions are fitted by a power law: $dN/dL \propto L^{-\alpha}$, where the normalization is expressed as the number of sources more luminous than $L_{\text{ref}} = 5 \times 10^{37}$ erg s$^{-1}$. Quoted errors for the slope $\alpha$ are at the 90% confidence level. In the case of Fornax-field sources, an additional, fixed component for the CXB has been included in the fit.

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According to Hou & Li (2016), the X-ray emission from UCDs is most likely due to dynamically formed LMXBs, as in the case of GCs. Here we conservatively distinguish the X-ray counterparts of UCDs and GCs.
about half of the excess sources (with \( L_X \gtrsim 10^{37} \text{ erg s}^{-1} \)) found in their sample of ETGs can be attributed to supernova-kicked LMXBs, which on average have an abundance of \( \sim 0.5 \) source per \( 10^{10} M_\odot \) (with \( L_X \gtrsim 5 \times 10^{37} \text{ erg s}^{-1} \), normalized to the host galaxy’s stellar mass). Assuming a similar scaling holds in Fornax, we estimate that \( \lesssim 34 \) supernova-kicked LMXBs are among our excess sources, a value obtained by summing up the 29 member galaxies including NGC 1399.

The third possibility is X-ray binaries associated with a more diffuse stellar population, either the extended stellar halo of NGC 1399 (Iodice et al. 2016) or the ICL (Iodice et al. 2017). The BCG halo grows primarily from mergers of infalling galaxies, whereas the ICL is built up mainly through stars tidally stripped from the outskirts of large galaxies or tidally disrupted low-mass galaxies orbiting in the cluster potential. Both being the result of hierarchical clustering, these two components are physically connected to each other, and it is noteworthy that a clear cut between them is rather difficult with morphological information only. Nevertheless, consisting of predominantly old stellar populations, both the BCG halo and ICL are expected to host LMXBs. In Section 4.1, we have attempted a spatial decomposition of the BCG halo contribution by fitting the observed surface density profile with a presumed exponential form, finding that this halo does appear to be the dominant component over the range of \( R \approx 3–8 R_e \) (Figure 4), potentially accounting for \( \sim 110 \) excess sources. In a recent study of X-ray sources in M87, the central giant elliptical galaxy of Virgo, Luan et al. (2018) found that the X-ray source surface density profile is consistent with the stellar halo distribution out to a projected radius of \( \sim 100 \) kpc. van Haften et al. (2018) also reported an excess of X-ray sources in the halo of NGC 4472, a group-central giant elliptical galaxy, although these authors essentially referred to the excess as sources without a known optical counterpart.

The exponential halo, however, cannot solely account for the detected excess. Subtracting the best-fitted halo contribution still leaves a residual of \( \sim 74 \) sources (or 40\% of all excess sources) spreading over a large radial range (Figure 4). This residual is likely related to the putative ICL. Based on the FDS, Iodice et al. (2017) found clear evidence for the presence of ICL in the core of Fornax, the bulk of which lies between NGC 1387, NGC 1379, and NGC 1381 (a region \( \sim 10’– 40’ \) west of NGC 1399) and possibly arises from tidal disruption of these three galaxies. Using the ICL g-band luminosity of \( 8.3 \times 10^8 L_\odot \) measured by Iodice et al. (2017), and adopting the values of \( M/L \) and \( c_\text{X} \) obtained in Section 4.1, we estimate that \( \sim 20 \) LMXBs could be associated with the ICL in this particular region. As a related remark, some GCs (and any associated GC-LMXBs) might also be classified as part of the ICL. Indeed, Iodice et al. (2017) found the aforementioned ICL to be spatially coincident with a previously known overdensity of blue GCs (D’Abrusco et al. 2016).

The last possible origin of the excess we consider, which is also closely related to hierarchical evolution in the cluster environment, is the relic of tidally stripped nucleated galaxies (Ferguson & Binggeli 1994). In particular, relic galactic nuclei can manifest themselves as UCDs (e.g., Liu et al. 2015), whose typical luminosities and sizes (hence stellar densities) are intermediate between the classical GCs and dwarf elliptical galaxies (Brüns & Kroupa 2012). Empirically \( \sim 3\% \) of the known UCDs have an X-ray counterpart with \( L_X \gtrsim 10^{37} \text{ erg s}^{-1} \) (Hou & Li 2016), which are most likely LMXBs formed via stellar encounters, as in the case of GCs. Alternatively, a massive BH, if existed at the center of the UCD, can also give rise to detectable X-rays. Such a case might be relevant to M60-UCD1 (Seth et al. 2014) and M59-UCD3 (Ahn et al. 2018), two very massive UCDs both showing strong evidence for a central massive BH. Among the 144 UCDs currently known in Fornax, we have identified 15 with an X-ray counterpart, but due to the somewhat ambiguous classification of UCDs, 13 of these have also been considered GCs (Section 3.1). In any case, UCDs and relic galactic nuclei are unlikely to have a substantial contribution to the excess sources. In passing, we note that no X-ray counterpart (with a 3\% upper limit of \( L_X < 2 \times 10^{37} \text{ erg s}^{-1} \)) is found for Fornax UCD3, which was recently suggested to host a central BH of 3.5-million solar masses (Afanasiev et al. 2018). A detailed study of the X-ray properties of the UCDs, in parallel with the identified GC-LMXBs in Fornax, will be presented elsewhere.

6. Concluding Remarks

Using extensive Chandra observations, we have studied the X-ray sources in the inner \( \sim 30’ \) (\( \sim 180 \) kpc) of the Fornax cluster. Our main results include:

1. We detect a total of 1279 independent X-ray sources, among which 1177 are detected in the 0.5–8 keV band down to a limiting luminosity of \( \sim 3 \times 10^{37} \text{ erg s}^{-1} \). By cross-correlating various optical catalogs, we identify 18 foreground sources, 270 GCs, 15 UCDs, and 4 nuclear sources. A source catalog is presented for future reference.

2. From the radial surface density profile, we statistically identify \( \sim 183 \) excess sources beyond three times the effective radius of NGC 1399, with respect to the expected CXB.

3. The luminosity function of the excess sources is significantly steeper than that of the GC-LMXBs, suggesting that any unidentified GCs would have only a minor contribution to the excess sources.

4. We find that LMXBs associated with either the extended stellar halo of NGC 1399 or the ICL can be responsible for the majority of the excess sources. Supernova-kicked LMXBs may account for a small but non-negligible fraction of the excess.

The above findings provide strong evidence for the presence of intracluster X-ray sources in Fornax, the second unambiguous case for a galaxy cluster after Virgo. Owing to the low surface brightness, it is a persistent challenge to detect ICL with conventional optical imaging. Discrete sources, such as GCs, planetary nebulae (Theuns & Warren 1997) and novae (Neill et al. 2005), have been proposed as a complementary tracer for the ICL. The detection of intracluster X-ray sources in Virgo (Hou et al. 2017) and Fornax (this work) now opens up a new window to study the ICL. The rather steep power-law slope of the excess sources in Fornax indicates that there could be many unresolved, fainter sources below the current detection limit, which can be probed with deeper Chandra observations. In the case of Virgo, where the large sky area renders it impractical to conduct a full mapping with Chandra or XMM-Newton, the upcoming all-sky survey by the eROSITA mission (Merloni et al. 2012) is expected to find at least 10 times more intracluster X-ray sources with \( L_X \gtrsim 2 \times 10^{38} \text{ erg s}^{-1} \). Such sources may form a statistically meaningful sample for
exploring the physical properties and formation history of the ICL in the two clusters.

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