Spin-Orbit Coupling Effect on the Electrophilicity Index, Chemical Potential, Hardness and Softness of Neutral Gold Clusters: A Relativistic Ab-initio Study

Mahnaz Jabbarzadeh Sani a

aDepartment of Chemistry, College of Science, Shiraz University, Shiraz, Iran.

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Abstract
Electrophilicity index (\(\omega\)) is related to the energy lowering associated with a maximum amount of electron flow between a donor and an acceptor and possesses adequate information regarding structure, stability, reactivity and interactions. Chemical potential (\(\mu\)) measures charge transfer from a system to another having a lower value of \(\mu\), while chemical hardness (\(\eta\)) is a measure of characterizing relative stability of clusters. The main purpose of the present research work is to examine the Spin-Orbit Coupling (SOC) effect on the behavior of the electrophilicity index, chemical potential, hardness and softness of neutral gold clusters \(\text{Au}_n\) \((n=2-6)\). Using the second-order Douglas-Kroll-Hess Hamiltonian, geometries are optimized at the DKH2-B3P86/DZP-DKH level of theory. Spin-orbit coupling energies are computed using the fourth-order Douglas-Kroll-Hess Hamiltonian, generalized Hartree-Fock method and all-electron relativistic double-\(\zeta\) level basis set. Then, spin-orbit coupling (SOC) corrections to the electrophilicity index, chemical potential, hardness and softness are calculated. It is revealed that spin-orbit correction to the vertical chemical hardness has important effect on \(\text{Au}_3\) and \(\text{Au}_6\), i.e. SOC decreases (increases) the hardness of gold trimer (hexamer). Due to the relationship between hardness and softness, \(\sigma = 1/\eta\), inclusion of spin-orbit coupling increases (decreases) the softness of \(\text{Au}_3\) (\(\text{Au}_6\)) and thus destabilizes (stabilizes) it. Spin-orbit coupling (SOC) also has more important effect on the chemical potential of \(\text{Au}_3\) by decreasing its value. It is found that spin-orbit coupling has considerable effect on the electrophilicity index of gold trimer and greatly increases its value. Furthermore, SOC increases the maximal charge acceptance of \(\text{Au}_3\) more and thus destabilizes it more. As a result, spin-orbit coupling effect appears to be important in calculating the electrophilicity index of the gold trimer.

Keywords: Density Functional Theory; Douglas-Kroll-Hess; Electronic Properties; Gold Clusters; Spin-orbit Coupling.

1. Introduction
As a kind of promising nanomaterials, metal nanoclusters (NCs) have sparked wide-spread attention. In recent years, gold nanoparticles (AuNPs) have been applied to biomedicine and biological sensing [1-5] due to their biocompatibility and unique physical properties. AuNPs are one of the most promising catalysts, in spite of bulk Au as an inactive material [6-8]. Most of the computations on small neutral gold clusters have been performed using spin-free (scalar-relativistic) methods [9-12]. To obtain reliable theoretical results for gold clusters, the scalar relativistic correction is substantial. However, the spin-orbit coupling is expected to be important. There exist theoretical studies regarding spin-orbit coupling effect using effective core potentials or plane-wave basis sets [13-19]. These studies mainly focused on the spin-orbit coupling effect on the highest-occupied lowest-unoccupied (HOMO-LUMO) energy gaps, geometries, binding energies per atom and optical absorption of gold clusters. Xiao and Wang [13] using

* Corresponding author: mjsani@shirazu.ac.ir

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PW91 functional and a plane-wave basis set concluded that the spin-orbit coupling increases the binding energy for all the clusters, Auₙ (n=5-11, 14, 20), within the same magnitude. They also found that spin-orbit coupling decreases the HOMO-LUMO energy gap, but has no effect on the magnetic moment. Using a spin-averaged relativistic effective core potential, Li et al. [14] evaluated the performance of various density functional theory methods on the geometries and atomization energies of Au₂, Au₃, Au₄ and Au₅. They included spin-orbit (SO) coupling effect using two-component spinor formalism. They found that the functionals including SO coupling effect overestimate the atomization energies of gold clusters compared with those just including the scalar-relativistic (SC) effect. Flores and Proupin [15] applying PBE functional and pseudopotentials have investigated the effect of the spin-orbit coupling on the structures, relative stability and HOMO-LUMO energy gap of the lowest-lying isomers of neutral Au₁₃ cluster. Rusakov et al. [16] have used two-component relativistic density functional theory with accurate small-core shape-consistent relativistic pseudopotentials and spin-orbit corrections to refine the isomerization energy profile of Au₁ computed by spin-orbit free coupled cluster methods. The noble metal trimers, Au₃, Ag₃ and Cu₃, have also been investigated applying four-component Kohn-Sham-Dirac equation in the framework of relativistic density functional theory [17].

The concept of electrophilicity has been known for several decades. Parr et al. [20] proposed a definition based on the energy lowering associated with a maximum amount of electron flow between two species. It has been revealed that electrophilicity possesses adequate information regarding structure, stability, reactivity and interactions [21]. In the present work, spin-orbit coupling (SOC) effects on the electronic properties of small neutral gold clusters, Auₙ (n=2–6), are investigated applying the Douglas-Kroll-Hess Hamiltonian and all-electron relativistic basis set of valence double-ζ quality plus polarization functions (DZP-DKH). The spin-orbit coupling effects on the electrophilicity index, electronic chemical potential, absolute hardness and softness are examined.

2. Computational Methods

All calculations are performed using the Gaussian 09 suite of program [22] and the plots of molecular configurations and contour maps are generated with the GaussView software [23]. The B3P86 functional is used in conjunction with the valence double-ζ quality plus polarization functions (DZP-DKH [8s7p4d2f]) basis set [24]. The B3P86 functional has already proven to perform well for ionization potential computations of small neutral gold clusters [10]. The second-order Douglas-Kroll-Hess Hamiltonian [25-29] is used instead of the Schrödinger operator. For geometry optimizations, the coordinates are chosen according to the experimentally determined structures by Gruene et al. [30]. Using the second-order Douglas-Kroll-Hess Hamiltonian, all geometries are fully optimized at the DHK2-B3P86/DZP-DKH level of theory followed by harmonic vibrational frequency analysis. Then, from these optimized geometries, the spin-orbit coupling (SOC) energies are calculated using the fourth-order Douglas-Kroll-Hess Hamiltonian and the generalized Hartree-Fock method [31]. The electronic energy including spin-orbit coupling, \( E_{SO} \), is calculated using the following definition,

\[
E_{SO} = E_{SC} + \Delta_{SO}E
\]  

Where, \( E_{SC} \) is the electronic energy obtained from spin-orbit free (scalar-relativistic) optimizations at the DHK2-B3P86/DZP-DKH level, and \( \Delta_{SO}E \) is the spin-orbit coupling energy calculated by the DHK2-B3P86/DZP-DKH //DKH/SO- GHF/DZP-DKH level of theory. In order to determine the spin-orbit coupling (SOC) effect on the electronic properties of neutral gold clusters, the spin-orbit corrections to the chemical hardness, softness, chemical potential and electrophilicity index [20, 32] of the gold clusters are computed. In the following, the spin-orbit correction to a particular property, \( \Delta_{SO}P \), is defined using,

\[
\Delta_{SO}P = P_{SO} - P_{SC}
\]

Where, \( P_{SO} \) and \( P_{SC} \) are the calculated properties with considering the spin-orbit coupling energy, \( E_{SO} \), and without considering the spin-orbit coupling energy, \( E_{SC} \), respectively. Within the valence state parabola model [20], the chemical potential, chemical hardness, softness and electrophilicity index are introduced by,

\[
\mu = \left( \frac{\partial E}{\partial N} \right)_V = \frac{IP + EA}{2}
\]

(3)

\[
\eta = \left( \frac{\partial^2 E}{\partial N^2} \right)_V = \frac{IP - EA}{2}
\]

(4)

\[
\sigma = \frac{\eta}{2}
\]

(5)

\[
\omega = \frac{\eta^2}{2n}
\]

(6)

Where \( N, V, IP \) and \( EA \) are the number of electrons, the potential due to the nuclei plus any external potential, ionization potential and electron affinity, respectively. The spin-orbit corrections to these properties are calculated from,
The aforementioned properties are calculated using the both vertical and adiabatic ionization potential and electron affinities. Figure 1 illustrates the steps in calculating the spin-orbit coupling (SOC) effect on the electrophilicity index, chemical potential, hardness and softness of neutral gold clusters.

1. Selection of the appropriate coordinates for the optimization.

2. Selection of the appropriate method and basis set for the geometry optimization (This work: DKH2-B3P86/DZP-DKH level).

3. Selection of the correct Hamiltonian for the relativistic calculations (This work: Douglas-Kroll-Hess (DKH) Hamiltonian).

4. Optimization of the selected geometries at the DKH2-B3P86/DZP-DKH level and applying DKH Hamiltonian.

5. Performing a harmonic vibrational frequency analysis to characterize the nature of the stationary points.

6. Calculation of the spin-orbit coupling energies at the DKH2-B3P86/DZP-DKH //DKHGO- GHF/DZP-DKH level of theory ($\Delta_{SO}E$ = \langle L, S \rangle = spin-orbit coupling energy).

7. Calculation of the electronic energy including spin-orbit coupling: $E_{SO} = E_{SC} + \Delta_{SO}E$

8. Calculation of the properties (using the both $E_{SC}$ and $E_{SO}$):
   - Chemical potential $\mu = \frac{IP+EA}{2}$
   - Hardness $\eta = \frac{IP- EA}{2}$
   - Softness $\sigma = \frac{1}{\eta}$
   - Electrophilicity index $\omega = \frac{\mu^2}{2\eta}$

9. Calculation of the spin-orbit corrections to these properties:
   - $\Delta_{SO}\mu = \mu_{SO} - \mu_{SC}$
   - $\Delta_{SO}\eta = \eta_{SO} - \eta_{SC}$
   - $\Delta_{SO}\sigma = \sigma_{SO} - \sigma_{SC}$
   - $\Delta_{SO}\omega = \omega_{SO} - \omega_{SC}$

Figure 1. Flow diagram showing the computational steps.
3. Results and Discussion

The optimized structures of gold clusters at the DKH2-B3P86/DZP-DKH level of theory are displayed in Figure 2. Table 1 compares the apex angle, symmetry and electronic state of the two Jahn-Teller components of gold trimer. The isomer with the lowest energy is an acute angle triangular structure with $C_{3v}$ symmetry, a 58.455° apex angle, and a 2.537 Å bond length. The result is in agreement with the coupled cluster calculations of Schwerdtfeger et al. [33]. In fact, because of the Jahn-Teller effect, the $D_{3h}$ symmetry of Au$_5$ distorts to an acute triangular $C_{3v}$ structure. The lowest-energy structures of Au$_4$ and Au$_5$ are a trapezoid with $C_{3v}$ symmetry and a W-shaped geometry, respectively. The $D_{3h}$ planar triangular structure of Au$_6$ is obtained by adding one gold atom to W-shaped planar Au$_5$.

Figure 2. The calculated structures of gold clusters (Au$_n$, n=3-6) at the DKH2-B3P86/DZP-DKH level of theory

Table 1. Apex angle, symmetry and electronic state of the two Jahn-Teller components of gold trimer

| Apex angle | Symmetry | Electronic state |
|------------|----------|------------------|
| Au$_3$ (acute) | 58.455° | $C_{3v}$ | $^3A_1$ |
| Au$_4$ (obtuse) | 62.447° | $C_{3v}$ | $^3B_2$ |

The calculated vertical and adiabatic ionization potential, electron affinity, chemical potential, hardness, softness and electrophilicity index values are given in Tables 2 to 5. Table 6 presents the spin-orbit coupling (SOC) corrections to the chemical potential, hardness, softness and electrophilicity index. Chemical hardness has been used to characterize the relative stability of clusters. The principle of maximum hardness (PMH) states that systems at equilibrium present the highest value of hardness [34]. The computed hardness values with and without considering spin-orbit coupling (SOC) effect are plotted in Figure 3a. According to this Figure, the chemical hardness computed using adiabatic electron affinities and ionization potentials exhibits an odd-even oscillation behavior, whether spin-orbit coupling is considered or not. The even-sized clusters with closed-shell electronic configurations have higher values of chemical hardness compared to their immediate open-shell neighbors, indicating their higher stability. This is in agreement with the previous study of Singh and Sarkar [11]. They performed spin-free (scalar- relativistic) calculations using B3LYP/LANL2DZ method. The spin-orbit correction to the chemical hardness with the increasing cluster size is presented in Figure 4a. As this Figure shows, the spin-orbit correction to the vertical chemical hardness has important effect on Au$_3$ and Au$_6$. The $\Delta_{\text{v,so}}\mu$ value for acute Au$_3$ and Au$_6$ is negative (-0.038 eV) and positive (0.045 eV), respectively, i.e. spin-orbit coupling decreases (increases) the hardness of gold trimer (hexamer).

The softness, $\sigma$, is simply the inverse of the hardness. Figure 3b depicts the variation of softness as a function of cluster size, with and without considering spin-orbit coupling (SOC) effect. As this Figure illustrates, the oscillation behavior of softness is opposite to that of chemical hardness, due to its relationship with hardness, $\sigma = 1/\eta$. Moreover, the obtuse Au$_3$ has the highest adiabatic softness ($\sigma_{a,3} = 0.454$ eV$^{-1}$ and $\sigma_{3,so} = 0.456$ eV$^{-1}$), indicating its high reactivity. On the other hand, hard molecules have a large HOMO-LUMO gap, whereas soft molecules have a small energy gap [32]. The HOMO-LUMO energy gap of Au$_2$, obtuse (acute) Au$_3$, Au$_4$, Au$_5$ and Au$_6$ at the DKH2-B3P86/DZP-DKH level of theory is 2.760, 1.170 (1.454), 1.982, 1.529 and 2.739 eV, respectively. Therefore, obtuse Au$_3$ with the highest adiabatic softness has the lowest HOMO-LUMO gap value of 1.170 eV. The variation of the spin-orbit correction to the softness of neutral gold clusters as a function of cluster size is plotted in Figure 4b. The spin-orbit correction to the softness shows an even-odd alternation behavior. Furthermore, the inclusion of spin-orbit coupling increases (decreases) the softness of Au$_3$ (Au$_6$) and thus destabilizes (stabilizes) it, $\Delta_{\text{v,so}}\sigma_{v,3} = +0.005$ eV ($\Delta_{\text{v,so}}\sigma_{v,6} = -0.006$ eV).

Chemical potential, $\mu$, is related to the charge transfer from a system to another having a lower value of $\mu$. Hence, it is anticipated that the odd-numbered Au$_n$ clusters present higher $\mu$ values because they have an open shell and that after transferring one electron, they will close their electronic shell and will be more stable than their original open-shell clusters. The variation of chemical potential with and without considering spin-orbit coupling (SOC) effect is depicted in Figure 3c. According to this Figure, the vertical chemical potential of obtuse Au$_3$ has the highest value, with or without considering spin-orbit coupling ($\mu_{v} = -4.712$ eV, $\mu_{v,3,so} = -4.750$ eV ), indicating its high spontaneous response and chemical reactivity.

Figure 4c shows the spin-orbit correction to the chemical potential of neutral gold clusters versus the cluster size. Spin-orbit coupling increases the vertical chemical potential of Au$_2$ by 0.021 eV, while decreases that of Au$_3$ and Au$_6$. 
by 0.038 and 0.037 eV, respectively. Hence, spin-orbit coupling has more important effect on the chemical potential of Au.

Electrophilicity has been a measure of the energy stabilization of a cluster due to acquiring additional electronic charge from its surroundings. It is expected that the electrophilicity index (\(\omega\)) should be related to the electron affinity (EA), because both \(\omega\) and EA measure the capability of an agent to accept electrons. However, EA reflects the capability of accepting only one electron from the environment, whereas the electrophilicity index (\(\omega\)) measures the energy lowering of a cluster due to maximal electron flow between donor and acceptor. The electron flows may be either less or more than one. Meanwhile, the electrophilicity index depends not only on EA, but also on IP. Electron affinity and electrophilicity index are related; yet they are not equal [20].

Table 2. Scalar-relativistic vertical ionization potential (\(v_{IP}/eV\)), vertical electron affinity (\(v_{EA}/eV\)), vertical chemical potential (\(\mu_{v}/eV\)), vertical hardness (\(\eta_{v}/eV\)), vertical softness (\(\sigma_{v}/eV^{-1}\)) and vertical electrophilicity index (\(\omega_{v}/eV\))

| n   | \(v_{IP}/eV\) | \(v_{EA}/eV\) | \(\mu_{v}/eV\) | \(\eta_{v}/eV\) | \(\sigma_{v}/eV^{-1}\) | \(\omega_{v}/eV\) |
|-----|----------------|----------------|----------------|----------------|----------------|----------------|
| 2   | 9.529          | 2.201          | -5.865         | 3.664          | 0.273          | 4.694          |
| 3   | (obtuse)       | 7.561          | 1.863          | -4.712         | 2.849          | 0.351          | 3.897          |
| 4   | 8.279          | 2.840          | -5.560         | 2.720          | 0.368          | 5.683          |
| 5   | 8.082          | 3.364          | -5.723         | 2.359          | 0.424          | 6.942          |
| 6   | 8.234          | 2.480          | -5.357         | 2.877          | 0.348          | 4.987          |

Table 3. Spin-orbit corrected vertical ionization potential (\(v_{IP_{so}}/eV\)), vertical electron affinity (\(v_{EA_{so}}/eV\)), vertical chemical potential (\(\mu_{v_{so}}/eV\)), vertical hardness (\(\eta_{v_{so}}/eV\)), vertical softness (\(\sigma_{v_{so}}/eV^{-1}\)) and vertical electrophilicity index (\(\omega_{v_{so}}/eV\))

| n   | \(v_{IP_{so}}/eV\) | \(v_{EA_{so}}/eV\) | \(\mu_{v_{so}}/eV\) | \(\eta_{v_{so}}/eV\) | \(\sigma_{v_{so}}/eV^{-1}\) | \(\omega_{v_{so}}/eV\) |
|-----|---------------------|---------------------|---------------------|---------------------|---------------------|---------------------|
| 2   | 9.489               | 2.199               | -5.844              | 3.645               | 0.274               | 4.685               |
| 3   | (obtuse)            | 7.561               | 1.938               | -4.750              | 2.812               | 0.356               | 4.012               |
| 4   | 8.286               | 2.837               | -5.562              | 2.725               | 0.367               | 5.676               |
| 5   | 8.079               | 3.368               | -5.724              | 2.356               | 0.424               | 6.953               |
| 6   | 8.316               | 2.472               | -5.394              | 2.922               | 0.342               | 4.979               |

Table 4. Scalar-relativistic adiabatic ionization potential (\(a_{IP}/eV\)), adiabatic electron affinity (\(a_{EA}/eV\)), adiabatic chemical potential (\(\mu_{a}/eV\)), adiabatic hardness (\(\eta_{a}/eV\)), adiabatic softness (\(\sigma_{a}/eV^{-1}\)) and adiabatic electrophilicity index (\(\omega_{a}/eV\))

| n   | \(a_{IP}/eV\) | \(a_{EA}/eV\) | \(\mu_{a}/eV\) | \(\eta_{a}/eV\) | \(\sigma_{a}/eV^{-1}\) | \(\omega_{a}/eV\) |
|-----|----------------|----------------|----------------|----------------|----------------|----------------|
| 2   | 9.523          | 2.221          | -5.872         | 3.651          | 0.274          | 4.722          |
| 3   | (obtuse)       | 7.547          | 3.145          | -5.346         | 2.201          | 0.454          | 6.492          |
| 4   | 8.240          | 2.878          | -5.559         | 2.674          | 0.374          | 4.440          |
| 5   | 8.090          | 3.381          | -5.695         | 2.314          | 0.432          | 7.008          |
| 6   | 8.182          | 2.528          | -5.355         | 2.827          | 0.354          | 5.072          |

Table 5. Spin-orbit corrected adiabatic ionization potential (\(a_{IP_{so}}/eV\)), adiabatic electron affinity (\(a_{EA_{so}}/eV\)), adiabatic chemical potential (\(\mu_{a_{so}}/eV\)), adiabatic hardness (\(\eta_{a_{so}}/eV\)), adiabatic softness (\(\sigma_{a_{so}}/eV^{-1}\)) and adiabatic electrophilicity index (\(\omega_{a_{so}}/eV\))

| n   | \(a_{IP_{so}}/eV\) | \(a_{EA_{so}}/eV\) | \(\mu_{a_{so}}/eV\) | \(\eta_{a_{so}}/eV\) | \(\sigma_{a_{so}}/eV^{-1}\) | \(\omega_{a_{so}}/eV\) |
|-----|---------------------|---------------------|---------------------|---------------------|---------------------|---------------------|
| 2   | 9.482               | 2.218               | -5.850              | 3.632               | 0.275               | 4.711               |
| 3   | (obtuse)            | 7.547               | 3.162               | -5.355              | 2.193               | 0.456               | 6.538               |
| 4   | 8.248               | 2.875               | -5.562              | 2.687               | 0.372               | 5.757               |
| 5   | 8.090               | 3.386               | -5.698              | 2.312               | 0.433               | 7.021               |
| 6   | 8.199               | 2.514               | -5.357              | 2.843               | 0.352               | 5.047               |
The variation of electrophilicity index as a function of cluster size, with and without considering spin-orbit coupling (SOC) effect is plotted in Figure 3d. As can be seen in this Figure, the electrophilicity index computed using adiabatic electron affinities and ionization potentials shows an odd-even oscillation behavior. The odd-numbered gold clusters present local maxima, and due to their open-shells, have more tendencies to accept electronic charge. On the other hand, the even-numbered and closed shell gold clusters are more stable and are less likely to acquire additional electronic charge. Hence, low electrophilicity index values are expected for them. However, using vertical IPs and EAs, the electrophilicity index behaves similar to vertical electron affinity (Table 2), and due to the linear geometry of the stable anionic gold trimer [33], acute Au3 has a very low vertical electrophilicity index (ωv = 3.894 eV) and the even-odd alternation is not observed (Figure 3d). To illustrate the spin-orbit coupling effect on the electrophilicity index of small neutral gold clusters, the variation of spin-orbit correction to the ω as a function of cluster size is presented in Figure 4d. An odd-even alternation behavior is obvious. Moreover, the spin-orbit correction to the vertical and adiabatic electrophilicity values of acute Au3 (Δωv = 0.116 eV and Δωa = 0.056 eV) and Au5 (Δωv = 0.011 eV and Δωa = 0.013 eV) are positive, indicating that spin-orbit coupling increases the electrophilicity index of these clusters. It is evident that spin-orbit coupling has significant effect on the vertical electrophilicity index of the gold trimer.

Maximal charge acceptance, ΔNmax = −μ/η, measures the maximal flow of electrons between a donor and an acceptor [20]. Figure 5a depicts the variation of the maximal charge acceptance of neutral gold clusters as a function of cluster size, with and without considering spin-orbit coupling. Similar to the electrophilicity index (Figure 3d), the adiabatic maximal charge acceptance exhibits odd-even alternation behavior, whether spin-orbit coupling is considered or not. Furthermore, the open-shell and odd-numbered Au5 and Au7 clusters have high adiabatic maximal charge acceptance values of 2.442 and 2.465, respectively, i.e. these clusters show high tendency to accept electrons.

### Table 6. Spin-orbit coupling (SOC) corrections to the vertical and adiabatic chemical potential (Δv/aμ/eV), softness (Δv/aσ/eV⁻¹) and electrophilicity index (Δv/aω/eV)

| n   | Δvμ/eV | Δvσ/eV | Δvω/eV | Δaμ/eV | Δaσ/eV | Δaω/eV |
|-----|--------|--------|--------|--------|--------|--------|
| 2   | 0.021  | 0.022  | -0.019 | -0.019 | 0.001  | 0.001  | -0.009 | -0.011 |
| 3  (obtuse) | -0.038  | -0.009 | -0.037 | -0.008 | 0.005  | 0.002  | 0.115  | 0.046  |
| 3  (acute) | -0.038  | -0.016 | -0.038 | -0.016 | 0.005  | 0.002  | 0.116  | 0.056  |
| 4   | -0.002 | -0.003 | 0.005  | 0.006  | -0.001 | -0.001 | -0.007 | -0.006 |
| 5   | -0.001 | -0.003 | -0.003 | -0.002 | 0.000  | 0.001  | 0.011  | 0.013  |
| 6   | -0.037 | -0.002 | 0.045  | 0.016  | -0.006 | -0.002 | -0.008 | -0.025 |

Subscripts ‘v’, ‘a’ and ‘so’ indicate ‘vertical’, ‘adiabatic’ and ‘spin-orbit’, respectively. Obtuse and acute stand for the two Jahn-Teller components of gold trimer.

![Figure 3. Variation of the (a) chemical hardness, (b) softness, (c) chemical potential and (d) electrophilicity index values with cluster size. Subscripts 'so' and 'v' indicate 'spin-orbit' and 'vertical', respectively.](image-url)
The spin-orbit coupling corrected adiabatic maximal charge acceptance per atom of neutral gold clusters are 0.805, 0.814 (0.613), 0.517, 0.493 and 0.314 atom\(^{-1}\) for Au\(_2\), obtuse (acute) Au\(_3\), Au\(_4\), Au\(_5\) and Au\(_6\), respectively. Obtuse Au\(_3\) has the highest adiabatic maximal charge acceptance per atom value of 0.814 and 0.810 atom\(^{-1}\) with and without considering spin-orbit coupling, respectively, indicating high reactivity of this cluster due to its active sites. On the other hand, Au\(_6\) has the lowest adiabatic maximal charge acceptance per atom value of 0.314 and 0.316 atom\(^{-1}\) with and without considering spin-orbit coupling, respectively, indicating its high stability. The results of Nguyen et al. [12] based on the second difference of energy and fragmentation energy calculations, also predicted high reactivity (stability) for Au\(_3\) (Au\(_6\)). As can be seen in Figure 5b, the spin-orbit correction to the maximal charge acceptance of neutral gold clusters exhibits an odd-even alternation behavior. Moreover, when the maximal charge acceptance is computed using vertical ionization potentials and electron affinities, inclusion of spin-orbit coupling increases the maximal charge acceptance of obtuse Au\(_3\) more (0.035 eV) and thus destabilizes it more. The natural population analysis (NPA) shows that the natural charge is positive (negative) for the apex atom of acute (obtuse) gold trimer, i.e. the natural charge changes its sign in going from acute (2A\(_1\)) to the obtuse (2B\(_1\)) Au\(_3\) [35, 36]. Figure 6 represents the natural charges and electrostatic potential (ESP) contour maps of the lowest-energy gold clusters Au\(_n\) (n = 3-6), at the DKH2-B3P86/DZP-DKH level. It is obvious that the apex atom of acute gold trimer has the highest natural charge value of +0.172, indicating its capability to accept additional electronic charge. Red lines in the electrostatic potential contour maps correspond to the negative charges, so these regions are responsible for the nucleophilic interactions.

Figure 4. Spin-orbit corrections to the (a) chemical hardness, (b) softness, (c) chemical potential and (d) electrophilicity index. Subscripts ‘so’ and ‘v’ indicate ‘spin-orbit’ and ‘vertical’, respectively.

Figure 5. Variation of (a) maximal charge acceptances with and without considering spin-orbit coupling, (b) spin-orbit corrections to the maximal charge acceptance. Subscripts ‘so’ and ‘v’ indicate ‘spin-orbit’ and ‘vertical’, respectively.
Figure 6. Electrostatic potential (ESP) contour maps (right panel) and natural charge values (left panel) of neutral gold clusters $\text{Au}_n$ ($n=3-6$)
Furthermore, the separations of angular momentum vectors. The variation of spin larger than Figure 7a depicts the variation of the SOC energies versus the cluster size. From this Figure, a linear relationship is based on ionization potentials IP = E[Au−] – E[Au+], electron affinities EA = E[Au+] – E[Au−], where E[Au+], E[Au−] and E[Au0] are the total energies of the neutral, cationic and anionic clusters, respectively. As a result, based on Equation 1, the spin-orbit corrections to the chemical potential, hardness, softness and electrophilicity index depend on the spin-orbit energies of the neutral, positively and negatively charged gold clusters, i.e. (L.S)+, (L.S)+ and (L.S)−. In the following, the behavior of spin-orbit coupling energy ΔSOE = (L.S) (Equation 1) is investigated. Tables 7 and 8 present the calculated spin-orbit coupling (SOC) energies for the neutral, cationic and anionic gold clusters. Figure 7a depicts the variation of the SOC energies versus the cluster size. From this Figure, a linear relationship is observed between the spin-orbit coupling energies and the cluster size. Furthermore, the values of SOC energies, (L.S)0, (L.S)+ and (L.S)−, vary in a relatively broad range, i.e. between -0.152 and -0.462 au. According to Tables 7 and 8, the spin-orbit coupling (SOC) energies of neutral open-shell and odd-numbered Au3 and Au5 with one unpaired electron are larger than those of positively and negatively charged closed shell gold trimer and pentamer, i.e. (L.S)0 is larger than (L.S)+ and (L.S)−. Hence, unpaired electron(s) can result in higher SOC energies. Moreover, the geometry of a cluster may influence the ΔSOE = (L.S) magnitude due to appropriate orientations of the total orbital and spin angular momentum vectors. The variation of spin-orbit coupling energies per atom, i.e. (L.S)0/n, (L.S)+/n and (L.S)−/n versus the cluster size is plotted in Figure 7b. From this Figure, considerable separations among the SOC energies per atom of the neutral, positively and negatively charged Au3, Au5, and Au7 clusters are observed. Furthermore, the separations of the spin-orbit coupling (SOC) energies per atom of the gold trimer are larger, resulting in higher SOC effects on the electronic properties of this cluster.

![Figure 7](image-url)

Table 7. Spin-orbit coupling energies of the neutral <L.S>0, cationic <L.S>+ and anionic <L.S>− gold clusters at the DKH2-B3P86/DZP-DKH//DKHSO-GHF/DZP-DKH level. The SOC energies of the cationic and anionic clusters are at the optimized geometries of neutral clusters.

| n  | <L.S>0 (a.u.) | <L.S>− (a.u.) | <L.S>+ (a.u.) | <L.S>0/n (a.u./atom) | <L.S>−/n (a.u./atom) | <L.S>+/n (a.u./atom) |
|----|--------------|--------------|--------------|----------------------|----------------------|----------------------|
| 2  | -0.15304     | -0.15451     | -0.15298     | -0.07652             | -0.07726             | -0.07649             |
| 3 (obtuse) | -0.22862     | -0.22946     | -0.23214     | -0.07621             | -0.07649             | -0.07738             |
| 3 (acute)  | -0.22943     | -0.22945     | -0.23219     | -0.07648             | -0.07648             | -0.07740             |
| 4   | -0.30670     | -0.30641     | -0.30659     | -0.07668             | -0.07660             | -0.07665             |
| 5   | -0.38385     | -0.38395     | -0.38400     | -0.07677             | -0.07679             | -0.07680             |
| 6   | -0.46151     | -0.45848     | -0.46120     | -0.07692             | -0.07641             | -0.07687             |

Table 8. Spin-orbit coupling energies of the neutral <L.S>0, cationic <L.S>+, and anionic <L.S>− gold clusters at the DKH2-B3P86/DZP-DKH//DKHSO-GHF/DZP-DKH level. <L.S>0, <L.S>−, and <L.S>− are the SOC energies at the optimized geometries of neutral, cationic and anionic gold clusters, respectively.

| n  | <L.S>0 (a.u.) | <L.S>− (a.u.) | <L.S>− (a.u.) | <L.S>0/n (a.u./atom) | <L.S>−/n (a.u./atom) | <L.S>−/n (a.u./atom) |
|----|--------------|--------------|--------------|----------------------|----------------------|----------------------|
| 2  | -0.15304     | -0.15455     | -0.15293     | -0.07652             | -0.07728             | -0.07647             |
| 3 (obtuse) | -0.22862     | -0.22945     | -0.23006     | -0.07621             | -0.07648             | -0.07669             |
| 3 (acute)  | -0.22943     | -0.22945     | -0.23060     | -0.07648             | -0.07648             | -0.07687             |
| 4   | -0.30670     | -0.30643     | -0.30661     | -0.07668             | -0.07661             | -0.07665             |
| 5   | -0.38385     | -0.38386     | -0.38403     | -0.07677             | -0.07677             | -0.07681             |
| 6   | -0.46151     | -0.46085     | -0.46101     | -0.07692             | -0.07681             | -0.07683             |
4. Conclusion

Gold clusters are attracting interest as the building blocks of novel nano-structural materials due to their biocompatibility and unique physical properties. Most of the computations on small neutral gold clusters have been performed using scalar-relativistic methods, however, the Spin-Orbit Coupling (SOC) effect is appeared to be important. Previous studies focused on spin-orbit coupling effects on HOMO-LUMO energy gaps, geometries, cohesive energies and optical absorption of gold clusters. Electrophilicity index is related to the energy lowering associated with a maximum amount of electron flow between a donor and an acceptor and provides information about structure, stability, reactivity and interactions. In the present research work, the behaviour of the electrophilicity index along with chemical potential, hardness and softness is investigated. The main purpose is to examine spin-orbit coupling (SOC) effects on these properties. We use generalized Hartree-Fock method in conjunction with the fourth-order Douglas-Kroll-Hess (DKH) Hamiltonian and all-electron relativistic double-$\zeta$-level basis set to compute the spin orbit coupling energies and corrections to the electrophilicity index, chemical potential, hardness and softness. It is revealed that spin-orbit correction to the vertical chemical hardness has important effect on Au$_3$ and Au$_6$ by decreasing (increasing) the hardness of gold trimer (hexamer). Moreover, spin-orbit correction to the softness exhibits an even-odd oscillation behaviour and the inclusion of SOC increases (decreases) the softness of Au$_3$ (Au$_6$) and thus destabilizes (stabilizes) it. Spin-orbit coupling has more important effect on the chemical potential of gold trimer by decreasing its value. The adiabatic electrophilicity index shows an odd-even alternation behaviour and odd-numbered gold clusters present local maxima due to their open-shells. The spin-orbit corrections to the electrophilicity index of acute Au$_3$ and Au$_6$ are positive, i.e. SO coupling increases the electrophilicity index of these clusters. Furthermore, the open-shell and odd-numbered gold trimer has the highest adiabatic maximal charge acceptance per atom, with and without considering spin-orbit coupling, indicating high reactivity of this cluster. In addition, spin-orbit coupling increases the maximal charge acceptance of Au$_3$ more and thus destabilizes it more. The natural population analysis reveals that the apex atom of the acute gold trimer has the highest natural charge, indicating its capability to accept additional electronic charges.

5. Supplementary Material

Supplementary material (Appendix I) for this article contains XYZ coordinates of the lowest-energy gold clusters optimized at the DKH2-B3P86/DZP-DKH level of theory.

6. Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Appendix I

XYZ coordinates of the small neutral gold clusters optimized at the DKH2-B3P86/DZP-DKH level of theory.

| Au₂   | Au₂ (obtuse) |
|-------|--------------|
| 0.000000 0.000000 1.190360 | 0.000000 1.239142 -0.738019 |
| 0.000000 0.000000 -1.190360 | 0.000000 -1.239142 -0.738019 |
|        | 0.000000 0.000000 1.476038 |

| Au₃ (acute) | Au₃ |
|-------------|-----|
| 0.000000 1.238856 -0.738055 | 0.000000 1.278972 0.000421 |
| 0.000000 -1.238856 -0.738055 | 2.193981 0.000000 -0.000421 |
| 0.000000 0.000000 1.476111 | 0.000000 -1.278972 0.000421 |

| Au₄   | Au₅   |
|-------|-------|
| -1.267923 1.352279 0.001733 | -0.732383 -1.360849 0.120640 |
| -2.496312 -0.862669 -0.039240 | 1.365676 2.537432 -0.115246 |
| 1.268335 1.352251 0.002184 | 1.514815 -2.451354 -0.118110 |
| -0.000125 -0.978835 0.074794 | -0.813203 1.315450 0.115580 |
| 2.496024 -0.863026 -0.039470 | 1.545674 0.046071 0.115273 |
|        | -2.880579 -0.086751 -0.118137 |