WEIGHTED STATIONARY PHASE OF HIGHER ORDERS

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Abstract. The subject matter of this paper is an integral with exponential oscillation of phase $f(x)$ weighted by $g(x)$ on a finite interval $[\alpha, \beta]$. When the phase $f(x)$ has a single stationary point in $(\alpha, \beta)$, an $n$th-order asymptotic expansion of this integral is proved for $n \geq 2$. This asymptotic expansion sharpens the classical result for $n = 1$ by M.N. Huxley. A similar asymptotic expansion was proved by Blomer, Khan and Young under the assumptions that $f(x)$ and $g(x)$ are smooth and $g(x)$ is compactly supported on $\mathbb{R}$. In the present paper, however, these functions are only assumed to be continuously differentiable on $[\alpha, \beta]2n+3$ and $2n+1$ times, respectively. Because there are no requirements on the vanishing of $g(x)$ and its derivatives at the endpoints $\alpha$ and $\beta$, the present asymptotic expansion contains explicit boundary terms in the main and error terms. The asymptotic expansion in this paper is thus applicable to a wider class of problems in analysis, analytic number theory and other fields.

1. Introduction

In this paper we will consider exponential integrals of the form

$$\int_{\alpha}^{\beta} g(x)e(f(x)) \, dx.$$  \hspace{1cm} (1.1)

When $f'(x)$ changes signs at a point $x = \gamma$ with $\alpha < \gamma < \beta$, Huxley [3] obtained a first-order asymptotic expansion of (1.1). This asymptotic expansion has been widely used as a standard technique in analytic number theory. This integral also plays an important role in harmonic analysis. In the case of $\alpha = -\infty$ and $\beta = \infty$, Wall [7], pp.38–39, proved an $n$th order asymptotic expansion of (1.1). Blomer, Khan and Young [1] reproved such an asymptotic expansion and computed the main terms.

What we will do in the present paper is to further refine the asymptotic expansion of (1.1) in two aspects. Firstly we will consider the case of finite lower and upper limits in (1.1) with $g(x)$ and its derivatives being not necessarily zeros at the endpoints of the integration interval, as in [3]. This will bring in boundary terms which will appear both in the main terms and the error terms. Detailed treatment of these boundary terms is lengthy, but they are necessary for a wider class of applications. Secondly the functions $f(x)$ and $g(x)$ will not be assumed to be $C^\infty$, as opposite to [7] and [1].

Now let us have an overview of the stationary phase expansion we will prove (Theorem 1.2):

$$\int_{\alpha}^{\beta} g(x)e(f(x)) \, dx = \frac{e(f(\gamma) \pm 1/8)}{\sqrt{|f''(\gamma)|}} \left(g(\gamma) + \sum_{j=1}^{n} \varpi_{2j} \frac{(-1)^j (2j - 1)!!}{(2\pi i f''(\gamma))^j} \right) + \text{Boundary terms + Error terms.}$$

Here $n$ is related to the smoothness of $f$ and $g$, $\gamma$ is the only zero of $f'(x)$ in $(\alpha, \beta)$, and $\varpi_{2j}$ are given by (1.11). Possible applications of our results include Salazar and Ye [6] on spectral square moments of

$$S_X(f; \alpha, \beta) = \sum_{n} \lambda_f(n)e(\alpha n^\beta)\phi\left(\frac{n}{X}\right)$$

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for $0 < \beta < 1$, $\alpha \in \mathbb{R}^\times$, $\phi \in C^\infty_c((1,2))$, and $f$ being a Maass form for $\Gamma_0(N)$, and McKee, Sun and Ye [5] on an improved subconvexity bound for a Rankin-Selberg $L$-function for $SL_2(\mathbb{Z})$ and $SL_3(\mathbb{Z})$ Maass forms.

Our first theorem is a weighted first derivative test, which strengthens Lemma 5.5.5 of [3], p.113, with more boundary terms and smaller error terms. Similar theorems have been proved and used by Jutila and Motohashi [4] (Lemma 6) and Blomer, Khan and Young [1] (Lemma 8.1). We will thus not give its proof here but note that our version is on a finite integration interval and comes with boundary terms. We will also need the specific form of error terms later.

**Theorem 1.1.** Let $f(x)$ be a real-valued function, $n+2$ times continuously differentiable for $\alpha \leq x \leq \beta$, and let $g(x)$ be a real-valued function, $n+1$ times continuously differentiable for $\alpha \leq x \leq \beta$. Suppose that there are positive parameters $M$, $N$, $T$, $U$, with $M \geq \beta - \alpha$, and positive constants $C_r$ such that for $\alpha \leq x \leq \beta$,

$$\left| f^{(r)}(x) \right| \leq C_r \frac{T}{M^r}, \quad \left| g^{(s)}(x) \right| \leq C_s \frac{U}{N^s},$$

for $r = 2, \ldots, n+2$, and $s = 0, \ldots, n+1$. If $f'(x)$ and $f''(x)$ do not change signs on the interval $[\alpha, \beta]$, we have

$$\int_{\alpha}^{\beta} g(x) e(f(x)) dx = \left[ e(f(x)) \sum_{i=1}^{n} H_i(x) \right]^{\beta}_\alpha + O\left( \frac{M}{N} \sum_{j=1}^{\lfloor n/2 \rfloor} \min_j \left| f^{(n+j+1)}_M \right| \sum_{t=j}^{n-j} \frac{1}{N^{n-j-t} M^t} \right)$$

$$+ O\left( \left( \frac{M}{N} + 1 \right) \frac{U}{N \min_j \left| f^{(n+j)} \right|} \right) + O\left( \sum_{j=1}^{n} \min_j \left| f^{(n+j+1)}_M \right| \sum_{t=0}^{n-j} \frac{1}{N^{n-j-t} M^t} \right),$$

where

$$H_1(x) = \frac{g(x)}{2\pi i f'(x)}, \quad H_i(x) = -\frac{H_{i-1}'(x)}{2\pi i f''(x)}$$

for $i = 2, \ldots, n$.

If $g(x) \equiv 1$ on $[\alpha, \beta]$, we may take $U = 1$ and $N$ arbitrarily large. Then the first two error terms in Theorem 1.1 are negligible, while in the third error term we may take only one term with $t = n - j$ in the inner sum. This way we can get an explicit first derivative test, which supersedes Lemma 5.5.1 of [3], p.104, with more boundary terms and smaller error terms.

Our next theorem is an $n$th-order asymptotic expansion of a weighted stationary phase integral.

**Theorem 1.2.** Let $f(x)$ be a real-valued function, $2n+3$ times continuously differentiable for $\alpha \leq x \leq \beta$, and $g(x)$ a real-valued function, $2n+1$ times continuously differentiable for $\alpha \leq x \leq \beta$. Let $H_k(x)$ be defined as in (1.2). Assume that there are positive parameters $M$, $N$, $T$, $U$ with

$$M \geq \beta - \alpha,$$

and positive constants $C_r$ such that for $\alpha \leq x \leq \beta$,

$$\left| f^{(r)}(x) \right| \leq C_r \frac{T}{M^r}, \quad \text{for } r = 2, \ldots, 2n+3,$$

$$f''(x) \geq \frac{T}{C_2 M^2}$$

and

$$\left| g^{(s)}(x) \right| \leq C_s \frac{U}{N^s}, \quad \text{for } s = 0, \ldots, 2n+1.$$
Suppose that \( f'(x) \) changes signs only at \( x = \gamma \), from negative to positive, with \( \alpha < \gamma < \beta \). Let

\[
\Delta = \min \left\{ \frac{\log 2}{C_2}, \frac{1}{\max_{2 \leq k \leq 2n+3} \{C_k\}} \right\}.
\]

If \( T \) is sufficiently large satisfying \( \Delta > 1 \), we have for \( n \geq 2 \) that

\[
\int_\alpha^\beta g(x)e(f(x))dx = \frac{e^x}{\sqrt{f''(\gamma)}} \left( g(\gamma) + \sum_{j=1}^n \omega_j (-1)^j (2j-1)!! \right) + \left[ e(f(x)) \sum_{i=1}^{n+1} H_i(x) \right]_{\alpha}^{\beta} + O(\sum_{j=1}^n \frac{1}{(\gamma - \alpha)^{n+2}} + \frac{1}{(\beta - \gamma)^{n+2}}) + O \left( \frac{UM^{2n+4}}{T^{n+2}N^{2n+1}} \left( \frac{1}{(\gamma - \alpha)^{2n+3}} + \frac{1}{(\beta - \gamma)^{2n+3}} \right) \right)
\]

where

\[
\lambda_k = \frac{f^{(k)}(\gamma)}{k!} \text{ for } k = 2, \ldots, 2n + 2,
\]

\[
\eta_k = \frac{g^{(k)}(\gamma)}{k!} \text{ for } k = 0, \ldots, 2n,
\]

and

\[
\omega_k = \eta_k + \sum_{\ell=0}^{k-1} \sum_{j=1}^{k-\ell} \frac{C_{k\ell j}}{\lambda_j^2} \sum_{3 \leq n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j},
\]

with \( C_{k\ell j} \) being some constant coefficients.

**Remark 1.** Theorem 1.2 also holds when \( f'(x) \) changes signs from positive to negative, by changing the sign of \( 1/8 \) and taking the absolute value of \( f''(\gamma) \) inside the radical sign on the right hand side of (1.8).

**Remark 2.** Blomer, Khan and Young’s Proposition 8.2 and Corollary 8.3 in [1] obtained the same main terms and the last big-\( O \) term as in (1.8), under the assumptions that \( f(x) \) and \( g(x) \) are smooth and \( g(x) \) is compactly supported on \( \mathbb{R} \). In Theorem 1.2 however, \( f(x) \) and \( g(x) \) are only assumed to be continuously differentiable on \([\alpha, \beta]\) \( 2n + 3 \) and \( 2n + 1 \) times, respectively. Because there are no requirements that \( g(x) \) and its derivatives vanish at the endpoints \( \alpha \) and \( \beta \), Theorem 1.2 is valid for a much wider class of functions \( f(x) \) and \( g(x) \). This is indeed the case in [6].

**Remark 3.** In [1] the parameters require a condition that \( N \leq \beta - \alpha \). In Theorem 1.2, 1.3 is assumed instead, which is the same as assumed in Huxley [3].

We end this introduction with an outline of the proof. In §2 we divide the integration interval into three parts: \([\alpha, u] \), \([u, v] \) and \([v, \beta] \). In the middle subinterval we change variables from \( x \) to \( y \) by \( f(x) - f(\gamma) = \lambda_2 y^2 \). The goal is to obtain a Taylor approximation (in \( y \)) for \( g(x) \frac{dx}{dy} \) and its \( y \)-derivatives. The Taylor approximation (without error) to \( g(x) \frac{dx}{dy} \) is given by \( \sum_{k=0}^{2n} \omega_k y^k \), where \( \omega_k \) is given by (1.11) above.

In §3, we estimate the main weighted stationary phase integral. The main term comes from integrating the middle range in \( y \). \( x \in [u, v] \) corresponds to \( y \in [-r, r] \) for \( r \) given by (2.4). This leads us to estimate integrals of the form

\[
\int_{-r}^r e^{(\lambda_2 y^2)} y^k dy,
\]
(for even $k$) which we do by an application of the probability integral

$$
\Phi(x) = \frac{1}{\sqrt{\pi}} \int_0^x e^{-t^2} dt,
$$

and estimates of this integral in Gradshteyn and Ryzhik [2]. We need to estimate the integral of the error

$$
Q(y) = g(x) \frac{dx}{dy} - \sum_{k=0}^{2n} \omega_k y^k,
$$

and its $y$ derivatives. This involves previous estimates and the second derivative test, found in Huxley [3]. Also used is a dyadic decomposition of the interval $[-r, r]$.

Our proof is different from those in [7] and [1]. The reason we need to cut $[\alpha, \beta]$ into three pieces is that our techniques can only be applied to a small neighborhood of $\gamma$. On outer subintervals $[\alpha, u]$ and $[v, \beta]$ we apply Theorem 1.1.

## 2. Lemmas for Theorem 1.2

Under the assumptions of Theorem 1.2, we have the following Taylor expansions at $x = \gamma$,

$$
f(x) = f(\gamma) + \sum_{k=2}^{2n+2} \lambda_k (x - \gamma)^k + \frac{f^{(2n+3)}(\eta_0)}{(2n+3)!} (x - \gamma)^{2n+3},
$$

and for $1 \leq i \leq 2n+2$

$$
f^{(i)}(x) = \sum_{k=\max\{2, i\}}^{2n+2} \lambda_k^{(i)} (x - \gamma)^{k-i} + \frac{f^{(2n+3)}(\eta_i)}{(2n+3-i)!} (x - \gamma)^{2n+3-i},
$$

where $\eta_0, \ldots, \eta_{2n+2}$ are numbers between $x$ and $\gamma$ depending on $x$. Here

$$
\lambda_k^{(i)} = \frac{f^{(k)}(\gamma)}{(k-i)!} = \frac{k!}{(k-i)!} \lambda_k
$$

for $\max\{2, i\} \leq k \leq 2n+2$.

Now we change variables from $x$ to $y = h(x - \gamma)$ by

$$
f(x) - f(\gamma) = \lambda_2 (h(x - \gamma))^2 = \lambda_2 y^2,
$$

such that $y = h(x - \gamma)$ has the same sign as that of $x - \gamma$. Define $h(\alpha - \gamma) = -r_1$ and $h(\beta - \gamma) = r_2$, i.e., $f(\alpha) = f(\gamma) + \lambda_2 r_1^2$ and $f(\beta) = f(\gamma) + \lambda_2 r_2^2$. We choose a number $r$ which satisfies that

$$
r = \min\{r_1, r_2, \Delta M\}.
$$

Since

$$
\lambda_2 r_1^2 = f(\alpha) - f(\gamma) = \frac{f''(\eta)}{2!} |\alpha - \gamma|^2
$$

for some $\eta \in (\alpha, \gamma)$, by (1.4) and (1.5) we have

$$
\frac{1}{C_2^2} \leq \frac{|\alpha - \gamma|^2}{r_1^2} = \frac{2\lambda_2}{f''(\eta)} \leq C_2^2.
$$

Hence we have

$$
\frac{r_1}{C_2} \leq |\alpha - \gamma| \leq C_2 r_1.
$$

Likewise

$$
\frac{r_2}{C_2} \leq |\beta - \gamma| \leq C_2 r_2.
$$

By (2.4), (2.5) and (2.6), we see that

$$
r \geq \min \left\{ \frac{|\alpha - \gamma|}{C_2}, \frac{|\beta - \gamma|}{C_2}, \frac{M}{T^{\frac{n+3}{n+4}}} \right\}
$$
Define \( u, v \) by \( h(u - \gamma) = -r \) and \( h(v - \gamma) = r \), i.e., \( f(u) = f(v) = f(\gamma) + \lambda_2 r^2 \). By (2.4) we see that \( \alpha \leq u < \gamma < v \leq \beta \). In this section, we will only consider \( x \) and \( y \) in

\[
- r \leq y \leq r, \quad u \leq x \leq v.
\]

By (2.1) we know that

\[
y^2 = (x - \gamma)^2 \left( 1 + \sum_{k=3}^{2n+2} \frac{\lambda_k}{\lambda_2} (x - \gamma)^{k-2} + \frac{f^{(2n+3)}(\theta_0)}{\lambda_2(2n+3)!} (x - \gamma)^{2n+1} \right).
\]

Similar to (2.5), we have

\[
(2.9) \quad \frac{r}{C_2} \leq |u - \gamma| \leq C_2 r \quad \text{and} \quad \frac{r}{C_2} \leq |v - \gamma| \leq C_2 r.
\]

By (2.4), (2.8) and (2.9), we see that

\[
|u| \leq \frac{r}{C_2} + 2.
\]

Note that by (1.7). By (2.10) we can get for \( 1 \leq \gamma \leq \beta \).

\[
\text{Consequently}
\]

\[
(2.11) \quad y^j = (x - \gamma)^j \left( 1 + \sum_{k=1}^{2n+2} \frac{\lambda_{k+2}}{\lambda_2} (x - \gamma)^{k} + \frac{f^{(2n+3)}(\theta_0)}{\lambda_2(2n+3)!} (x - \gamma)^{2n+1} \right)^{j/2}
\]

by Taylor expansion at \( x = \gamma \). Here the Taylor coefficients \( \mu_{jk} \) can be determined by applying binomial expansions to (2.11)

\[
y^j = (x - \gamma)^j \left( 1 + \sum_{i=1}^{\infty} C_{ji} \left( \sum_{k=1}^{2n} \frac{\lambda_{k+2}}{\lambda_2} (x - \gamma)^{k} + \frac{f^{(2n+3)}(\theta_0)}{\lambda_2(2n+3)!} (x - \gamma)^{2n+1} \right)^i \right)
\]

with

\[
C_{ji} = \binom{j/2}{i} \prod_{m=0}^{i-1} \left( \frac{j}{2} - m \right).
\]

Consequently

\[
(2.12) \quad \mu_{j0} = 1, \quad \mu_{jk} = \sum_{i=1}^{k} \frac{C_{ji}}{\lambda_2} \sum_{3 \leq n_1, \ldots, n_i \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_i} \quad \text{for} \quad 1 \leq k \leq 2n.
\]

The variable change between \( x \) and \( y \) in (2.3) and (2.11) allows us to express \( f^{(i)}(x) \) in terms of \( y \) for \( 1 \leq i \leq 2n + 2 \).
Lemma 2.1. Suppose \((1.4)\) and \((1.5)\) hold for \(f(x)\). For \(x\) and \(y\) in \((2.8)\) with \(r\) in \((2.4)\) we have

\[
f'(x) = \sum_{k=1}^{2n+1} \theta_k^{(1)} y^k + O_n \left( \frac{T|y|^{2n+2}}{M^{2n+3}} \right),
\]

where

\[
\theta_1^{(1)} = \lambda_2^{(1)} = 2\lambda_2, \quad \theta_k^{(1)} = \sum_{j=1}^{k-1} \frac{C_{k,j}}{\lambda_2^j} \sum_{3 \leq n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j} \quad \text{for} \quad 2 \leq k \leq 2n+1.
\]

Proof. We claim that for \(1 \leq m \leq 2n+1\)

\[
f'(x) = \sum_{k=1}^{m} \theta_{m,k} y^k + \sum_{k=m+1}^{2n+1} \theta_{m,k}(x-\gamma)^k + O_n \left( \frac{T|x-\gamma|^{2n+2}}{M^{2n+3}} \right),
\]

where

\[
\theta_{m,1} = \lambda_2^{(1)} = 2\lambda_2, \quad \theta_{m,k} = \sum_{1 \leq j \leq k-1} \frac{C_{k,j}}{\lambda_2^j} \sum_{n_1 + \cdots + n_j = k-1 + 2j} \lambda_{n_1} \cdots \lambda_{n_j} \quad \text{for} \quad 2 \leq k \leq 2n+1,
\]

which can be proved by induction. Take \(m = 2n+1\) in \((2.13)\) we get

\[
f'(x) = \sum_{k=1}^{2n+1} \theta_{2n+1,k} y^k + O_n \left( \frac{T|x-\gamma|^{2n+2}}{M^{2n+3}} \right) = \sum_{k=1}^{2n+1} \theta_k^{(1)} y^k + O_n \left( \frac{T|x-\gamma|^{2n+2}}{M^{2n+3}} \right).
\]

Using \((2.3)\) and from the second order Taylor expansion we see that

\[
\lambda_2 y^2 = f(x) - f(\gamma) = \frac{f''(w)}{2!} (x-\gamma)^2,
\]

where \(w\) is some constant between \(x\) and \(\gamma\). Then by \((1.4)\) and \((1.5)\)

\[
\frac{1}{C_2^2} \leq \frac{|x-\gamma|^2}{y^2} = \frac{2\lambda_2}{f''(\eta)} \leq C_2^2.
\]

Hence similar to \((2.9)\) we get \(|x-\gamma|/C_2 \leq y \leq C_2|x-\gamma|\). Then using above estimates we get

\[
f'(x) = \sum_{k=1}^{2n+1} \theta_k^{(1)} y^k + O_n \left( \frac{T|y|^{2n+2}}{M^{2n+3}} \right).
\]

\[
\square
\]

Similarly, we can change \(x\) to \(y\) in \((2.2)\) by \((2.11)\). We have for \(2 \leq i \leq 2n+2\)

\[
f^{(i)}(x) = \sum_{k=0}^{2n+2-i} \theta_{k}^{(i)} y^k + O_n \left( \frac{T|y|^{2n+3-i}}{M^{2n+3}} \right),
\]

where

\[
\theta_0^{(i)} = \lambda_i^{(i)} = i! \lambda_i, \quad \theta_k^{(i)} = \sum_{j=1}^{k} \frac{C_{k,j}}{\lambda_2^j} \sum_{3 \leq n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j} \quad \text{for} \quad 1 \leq k \leq 2n+2-i.
\]

Now by the definition of \(y\) in \((2.3)\) we can compute \(\frac{dx}{dy}\)

\[
\frac{dx}{dy} = \sum_{k=0}^{2n} \theta_k y^k + O_n \left( \frac{|y|^{2n+1}}{M^{2n+1}} \right).
\]

Lemma 2.2. With the above notation we assume \((1.4)\) and \((1.5)\). Then

\[
\frac{dx}{dy} = \sum_{k=0}^{2n} \theta_k y^k + O_n \left( \frac{|y|^{2n+1}}{M^{2n+1}} \right).
\]
where

\[ \rho_0 = 1, \quad \rho_k = \sum_{j=1}^{k} \frac{C_{kj}}{\lambda_j^2}, \quad \sum_{3 \leq n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j} \text{ for } k \geq 1. \]

**Proof.** First by (2.3), we see that

\[ \frac{dx}{dy} = \frac{2\lambda_2 y}{f'(x)}. \]

Then by (2.2) with \( i = 1 \) we get

\[ \frac{dx}{dy} = \frac{2\lambda_2 y}{2\lambda_2(x - \gamma)} \left( 1 + \sum_{k=3}^{2n+2} \frac{k\lambda_k}{2\lambda_2} (x - \gamma)^{k-2} + \frac{f'(2n+3)(\eta)}{2\lambda_2(2n+2)!} (x - \gamma)^{2n+1} \right). \]

Similar to (2.10), we can prove that

\[ \left| \sum_{k=3}^{2n+2} \frac{k\lambda_k}{2\lambda_2} (x - \gamma)^{k-2} + \frac{f'(2n+3)(\eta)}{2\lambda_2(2n+2)!} (x - \gamma)^{2n+1} \right| \]

\[ \leq \sum_{k=3}^{2n+2} \frac{C_kC_2(C_2\Delta)^{k-2}}{(k-1)!} < \frac{C_2^3 \max\{C_k\Delta\}}{2} \sum_{k=0}^{n} \frac{(C_2\Delta)^k}{k!} \]

\[ < \frac{C_2^3 \max\{C_k\Delta\} e^{C_2\Delta}}{2} \leq C_2^3 \max\{C_k\Delta\} \leq 1 \]

because of (2.10). Therefore, we get

\[ \frac{dx}{dy} = \frac{y}{x - \gamma} - \sum_{j=0}^{\infty} (-1)^j \left( \sum_{k=3}^{2n+2} \frac{k\lambda_k}{2\lambda_2} (x - \gamma)^{k-2} + \frac{f'(2n+3)(\eta)}{2\lambda_2(2n+2)!} (x - \gamma)^{2n+1} \right)^j \]

\[ = \frac{y}{x - \gamma} \left( 1 + \sum_{k=1}^{2n} \mu'_k(x - \gamma)^k + O_n\left( \frac{|x - \gamma|^{2n+1}}{M^{2n+1}} \right) \right) \]

with

\[ \mu'_0 = 1, \quad \mu'_k = \sum_{i=1}^{k} \frac{C_{i'}}{\lambda_{i'}^2}, \quad \sum_{3 \leq n_1, \ldots, n_i \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_i} \text{ for } 1 \leq k \leq 2n. \]

Now by (2.11) with \( j = 1 \), we have (noting (2.10))

\[ \frac{y}{x - \gamma} = \left( 1 + \sum_{k=1}^{2n} \frac{\lambda_{k+2}}{\lambda_2} (x - \gamma)^k + \frac{f'(2n+3)(\eta)}{\lambda_2(2n+3)!} (x - \gamma)^{2n+1} \right)^{1/2} \]

\[ = 1 + \sum_{k=1}^{2n} \mu''_k(x - \gamma)^k + O_n\left( \frac{|x - \gamma|^{2n+1}}{M^{2n+1}} \right) \]

with

\[ \mu''_0 = 1, \quad \mu''_k = \sum_{i=1}^{k} \frac{C_{i'}}{\lambda_{i'}^2}, \quad \sum_{3 \leq n_1, \ldots, n_i \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_i} \text{ for } 1 \leq k \leq 2n. \]

Then by (2.18) and (2.19), we conclude that

\[ \frac{dx}{dy} = \left( 1 + \sum_{k=1}^{2n} \mu'_k(x - \gamma)^k + O_n\left( \frac{|x - \gamma|^{2n+1}}{M^{2n+1}} \right) \right) \left( 1 + \sum_{k=1}^{2n} \mu''_k(x - \gamma)^k + O_n\left( \frac{|x - \gamma|^{2n+1}}{M^{2n+1}} \right) \right) \]

\[ \frac{dx}{dy} = \sum_{k=0}^{2n} \mu''_k(x - \gamma)^k + O_n\left( \frac{|x - \gamma|^{2n+1}}{M^{2n+1}} \right). \]
with $\mu'''_k = \sum_{i=0}^{k} \mu'_i \mu''_{k-i}$. Since (ignoring all coefficients)

$$
\mu'_k \mu''_l = \sum_{1 \leq s \leq k} \frac{1}{\lambda^s_2} \sum_{3 \leq n_i \leq 2n+3} \sum_{n_1 + \cdots + n_s = k + 2s} \lambda_{n_1} \cdots \lambda_{n_s}
\times \sum_{1 \leq s \leq l} \frac{1}{\lambda^s_2} \sum_{3 \leq n_i \leq 2n+3} \sum_{n_1 + \cdots + n_s = l + 2s} \lambda_{n_1} \cdots \lambda_{n_s}
= \sum_{1 \leq s \leq k+l} \frac{1}{\lambda^s_2} \sum_{3 \leq n_i \leq 2n+3} \sum_{n_1 + \cdots + n_s = k + l + 2s} \lambda_{n_1} \cdots \lambda_{n_s},
$$

we get that

$$
\mu'''_0 = 1, \mu'''_k = \sum_{i=1}^{k} \frac{C'''_{ki}}{\lambda^i_2} \sum_{3 \leq n_i, \ldots, n_i \leq 2n+3} \sum_{n_1 + \cdots + n_i = k + 2i} \lambda_{n_1} \cdots \lambda_{n_i}, \text{ for } 1 \leq k \leq 2n.
$$

Finally, changing $x$ to $y$ in (2.20), we get

$$
\frac{dx}{dy} = \sum_{k=0}^{2n} \rho_k y^k + O_n \left( \frac{|y|^{2n+1}}{M^{2n+1}} \right)
$$

where

$$
\rho_0 = 1, \rho_k = \sum_{j=1}^{k} \frac{C'''_{kj}}{\lambda^j_2} \sum_{3 \leq n_i, \ldots, n_j \leq 2n+3} \sum_{n_1 + \cdots + n_j = k + 2j} \lambda_{n_1} \cdots \lambda_{n_j}, \text{ for } k \geq 1,
$$

which can be proved by induction like Lemma 2.21. □

By the assumptions of Theorem 1.2, we also have the following Taylor expansions

$$(2.21)\quad g(x) = \sum_{k=0}^{2n} \eta_k (x-\gamma)^k + O_n \left( \frac{|x-\gamma|^{2n+1}}{N^{2n+1}} \right),$$

and

$$(2.22)\quad \frac{d^i g}{dx^i} = \sum_{k=0}^{2n} \eta^{(i)}_k (x-\gamma)^{k-i} + O_n \left( \frac{|x-\gamma|^{2n+1-i}}{N^{2n+1}} \right),$$

with

$$
\eta^{(i)}_k = \frac{g^{(k+i)}(\gamma)}{(k+i)!} = \frac{k!}{(k-i)!} \eta_k \text{ for } i \leq k \leq 2n.
$$

Similarly, if we change variables in (2.21) to $y$, we can get

$$(2.23)\quad g(x) = \sum_{k=0}^{2n} \eta'_k y^k + O_n \left( \frac{|y|^{2n+1}}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right), \text{ with } \eta'_0 = g(\gamma).$$

To determine other $\eta'_k$, we substitute (2.11) into (2.23) to get

$$
g(x) = \eta'_0 + \sum_{k=1}^{2n} \eta'_k (x-\gamma)^k \sum_{\ell=0}^{2n} \mu_{k\ell} (x-\gamma)^\ell + O_n \left( \frac{|x-\gamma|^{2n+1}}{M^{2n+1}} \right) + O_n \left( \frac{|y|^{2n+1}}{N^{2n+1}} \right)
= \eta'_0 + \sum_{m=1}^{2n} (x-\gamma)^m \sum_{k \geq 1, \ell \geq 0, k + \ell = m} \eta'_k \mu_{k\ell} + O_n \left( \frac{|y|^{2n+1}}{N^{2n+1}} \right).
$$
Consequently

(2.24) \[ \eta_1 = \sum_{k \geq 1, \ell \geq 0} \eta_k^1 \mu_{k\ell} = \eta_1^0 = \eta_1', \]

and for \(2 \leq m \leq 2n\)

(2.25) \[ \eta_m = \sum_{k \geq 1, \ell \geq 0} \eta_k^m \mu_{k\ell} = \eta_m' + \sum_{k, \ell \geq 1} \eta_k^m \mu_{k\ell}, \]

where \(\eta_m\) is defined in (2.21). We may compute \(\eta_m'\) for \(m \geq 2\) recursively using (2.25).

Lemma 2.3. With the above notation

(2.26) \[ \eta_m' = \eta_m + \sum_{k=1}^{m-1} \eta_k \sum_{j=1}^{m-k} C_{mkj} \lambda_j \]

where the \(k\) sum vanishes when \(m \leq 1\).

Proof. By (2.24), (2.26) holds for \(m = 1\). Suppose (2.26) holds for any number \(\leq m\). Then by (2.25) and (2.12)

(2.27) \[ \eta_{m+1}' = \eta_{m+1} - \sum_{\ell=1}^{m} \eta_{m+1-\ell} \sum_{i=1}^{m+\ell} \sum_{m_1, \ldots, m_i \leq 2n+3} \lambda_{m_1} \cdots \lambda_{m_i} \]

The first two terms on the right side of (2.27) fit (2.26) for \(m + 1\). For the third term, we change the order of sums on \(\ell\) and \(k\), let \(h = i + j\), denote \(m_1, \ldots, m_i, n_{i1}, \ldots, n_j\) by \(p_1, \ldots, p_h\), and get

\[ -\sum_{k=1}^{m-1} \eta_k \sum_{\ell=k+1}^{m} \sum_{h=2}^{m+1-k} \frac{1}{\lambda_h^2} \sum_{i, j \geq 1, i + j = h} C_{\ell i} C_{\ell k j} \sum_{3 \leq p_1, \ldots, p_h \leq 2n+3} \lambda_{p_1} \cdots \lambda_{p_h} \]

which also fits (2.26) for \(m + 1\). \(\square\)

Similarly using (2.11) in (2.22) we get for \(1 \leq i \leq n + 1\)

(2.28) \[ \frac{d^i g}{dx^i} = \sum_{k=0}^{2n-i} \eta_k^{(i)} y^k + O_n \left( |y|^{2n+1-i} \left( \frac{1}{NM^{2n}} + \frac{1}{N^{2n+1}} \right) \right), \]

where

(2.29) \[ \eta_k^{(i)} = \frac{(k + i)!}{k!} \eta_{k+i} + \sum_{m=1}^{k+i-1} \eta_m \sum_{j=1}^{k+i-m} C_{kmj} \lambda_j^3 \sum_{3 \leq n_{i1}, \ldots, n_{ij} \leq 2n+3} \lambda_{n_{i1}} \cdots \lambda_{n_{ij}}. \]
Multiplying \( g(x) \) in (2.23) with \( \frac{dx}{dy} \) in (2.15) and using

\[
\eta_k^{(i)} \ll \left( \frac{U}{N^k} + \frac{U}{NM^{k-1}} \right), \quad \rho_k \ll \frac{1}{M^k},
\]

we get

\[
g(x) \frac{dx}{dy} = \sum_{k=0}^{2n} \varpi_k y^k + O_N \left( U |y|^{2n+1} \left( \frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right) \right),
\]

where

\[
\varpi_k = \sum_{\ell=0}^{k} \eta_{\ell}^k \rho_{k-\ell}.
\]

Note that by (2.16) and (2.23), \( \varpi_0 = \eta_0^0 \rho_0 = g(\gamma) \).

**Lemma 2.4.** With the notation as above, (1.11) holds.

**Proof.** By (2.31), (2.26) and (2.16) we have

\[
\varpi_N = \sum_{m=0}^{N-1} \left( \sum_{k=1}^{m-1} \sum_{j=1}^{m-k} \frac{C_{mkj}}{\lambda_2^j} \sum_{n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j} \right)
\]

\[
\times \sum_{h=1}^{N-m} \frac{C_{N-h}}{\lambda_2^h} \sum_{n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j},
\]

where the first \( k \) sum vanishes when \( m \leq 1 \). Note that (2.33) is contained in (1.11), while \( \sum_{m=0}^{N-1} \eta_m \) times (2.32) is also contained in (1.11). The rest sum of \( \varpi_N \) equals

\[
\sum_{m=2}^{N-1} \sum_{k=1}^{m-1} \sum_{j=1}^{m-k} \frac{C_{mkj}}{\lambda_2^j} \sum_{n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j}
\]

\[
\times \sum_{h=1}^{N-m} \frac{C_{N-h}}{\lambda_2^h} \sum_{n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j}
\]

\[
= \sum_{k=1}^{N-2} \eta_k \sum_{m=k+1}^{N-1} \sum_{i=1}^{N-k} \frac{1}{\lambda_2^i} \sum_{h, j \geq 1}^{m \leq h + j \leq N - m + 2h} \sum_{n_1, \ldots, n_j \leq 2n+3} \lambda_{n_1} \cdots \lambda_{n_j},
\]

which also fits (1.11). \( \square \)

In the following lemma we compute derivatives of \( g(x) \frac{dx}{dy} \), which we will use to prove our main theorem.

**Lemma 2.5.** With the above notation we assume (1.4) and (1.5). Then for \( 1 \leq i \leq n+2 \)

\[
\frac{d^i}{dy^i} \left( g(x) \frac{dx}{dy} \right) = \sum_{k=0}^{2n-i} \frac{(k+i)!}{k!} \varpi_{k+i} y^k + O_N \left( U |y|^{2n+i-1} \left( \frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right) \right).
\]
Proof. By (1.11) we see that the expression of \( \varpi_k \), \( 0 \leq k \leq 2n \), only uses \( \eta_k \), \( 0 \leq \ell \leq 2n \), and hence it only uses \( g^{(\ell)}(\gamma) \) for \( 0 \leq \ell \leq 2n \) by (1.10). By the same (1.11), \( \varpi_k \), \( 0 \leq k \leq 2n \), only requires \( \lambda_{n_1}, \ldots, \lambda_{n_j} \) for \( n_1, \ldots, n_j \leq 2n + 2 \). Thus by (1.9), it only requires \( f^{(\ell)}(\gamma) \), \( \ell = 2, \ldots, 2n + 2 \). Consequently, \( \varpi_k \) for \( 0 \leq k \leq 2n \) are independent of \( y \), and the terms \( \sum_{k=0}^{2n} \varpi_k y^k \) are the corresponding terms in the Taylor expansion of \( g(x) \frac{dx}{dy} \).

This implies that

\[
(2.35) \quad \frac{d^i}{dy^i} \left(g(x) \frac{dx}{dy}\right) = \sum_{k=0}^{2n-i} \frac{(k+i)!}{k!} \varpi_{k+i} y^k + R_i(y).
\]

where \( R_i(y) \) is the remainder term. We want to show

\[
(2.36) \quad R_i(y) \ll O_n \left(U|y|^{2n+1-i} \left(\frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}}\right)\right).
\]

In the following we will only consider the case of \( i = 1 \). Other cases are similar.

From (2.28) and (2.29)

\[
(2.37) \quad \frac{dg}{dx} = \sum_{k=0}^{2n-1} \eta_k^{(1)'} y^k + O_n \left(U|y|^{2n} \left(\frac{1}{NM^{2n}} + \frac{1}{N^{2n+1}}\right)\right),
\]

where

\[
\eta_k^{(1)'} = \frac{(k+1)!}{k!} \eta_{k+1} + \sum_{m=1}^{k} \sum_{j=1}^{k+1-m} \frac{C_{kmj}}{\lambda_j^2} \lambda_{n_1} \cdots \lambda_{n_j}
\]

By (2.17) we see that

\[
\frac{d^2x}{dy^2} = 2\lambda_2 x f'' - 2\lambda_2 y f'' \frac{dx}{(f')^2} = \frac{1}{y^3} \left(dx \frac{dx}{dy} - \frac{1}{2\lambda_2} \left(dx \frac{dx}{dy}\right)^3 f''\right).
\]

From last equation, (2.14) for \( i = 2 \), and (2.15) we see that \( \frac{d^2x}{dy^2} \) can be expressed as power series of \( y \), i.e.

\[
(2.38) \quad \frac{d^2x}{dy^2} = \sum_{k=0}^{2n-1} \rho_k^{(1)} y^k + O\left(\frac{|y|^{2n}}{M^{2n+1}}\right),
\]

where

\[
\rho_k^{(1)} = \sum_{j=1}^{k+1} \frac{C_{kj}}{\lambda_j^2} \lambda_{n_1} \cdots \lambda_{n_j}.
\]

With these preparations, we compute

\[
\frac{d}{dy} \left(g(x) \frac{dx}{dy}\right) = \frac{dg}{dx} \left(\frac{dx}{dy}\right)^2 + g(x) \frac{d^2x}{dy^2}.
\]

By (2.37), (2.15), (2.23) and (2.38) we get

\[
(2.39) \quad \frac{d}{dy} \left(g(x) \frac{dx}{dy}\right) = \sum_{m=0}^{2n-1} \varpi_m^{(1)} y^m + O_n \left(U|y|^{2n} \left(\frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}}\right)\right)
\]

with

\[
\varpi_m^{(1)} = \sum_{k, \ell, \xi \geq 0 \atop k + \xi + \ell = m} \eta_k^{(1)'} \rho_{\ell} \rho_{m-\ell} + \sum_{k, \ell \geq 0 \atop k + \ell = m} \eta_k^{(1)'} \rho_{\ell}^{(1)}.
\]

For \( 0 \leq m \leq 2n-1 \), \( \varpi_m^{(1)} \) as above involves \( \eta_k \), \( 0 \leq \ell \leq 2n \), and \( \lambda_{n_1}, \ldots, \lambda_{n_j} \) with \( 3 \leq n_1, \ldots, n_j \leq 2n + 2 \), and hence is independent of \( y \). Consequently, the terms for \( 0 \leq m \leq 2n-1 \) in (2.39) are terms in the Taylor expansion of \( \frac{d}{dy} \left(g(x) \frac{dx}{dy}\right) \). Comparing this with the Taylor terms in (2.35) for \( i = 1 \), we conclude that for
0 \leq k \leq 2n - 1 we have \((k + 1)\varpi_{k+1} = \varpi_k^{(1)}\) by the uniqueness of Taylor expansions. Therefore we see that (2.36) holds for \(i = 1\).

3. PROOF OF THEOREM 1.2

Recall that (2.1), we have for \(\alpha \leq x \leq u\)

\[ |f'(x)| \geq \frac{T|x - \gamma|}{C_2 M^2} \geq \frac{T|u - \gamma|}{C_2 M^2} \geq \frac{T}{C_2 M^2}. \]

Therefore by Theorem 1.1 for \(n + 1\) we get

\[
\int_\alpha^u g(x)e(f(x))dx + \int_v^\beta g(x)e(f(x))dx = \left[ e(f(x)) \sum_{i=1}^{n+1} H_i(x) \right]_\alpha^u + \left[ e(f(x)) \sum_{i=1}^n H_i(x) \right]_v^\beta + O \left( \frac{UM^{2n+5}}{T^{n+2}r^{n+2}} \sum_{j=1}^{\frac{n+1}{r}} \sum_{t=0}^{n+1-j} \frac{1}{N^{n+1-j-t}M^t} \right) + O \left( \frac{UM^{2n+4}}{T^{n+2}r^{n+2}} \sum_{j=1}^{\frac{n+1}{r}} \left( \frac{1}{(\gamma - \alpha)^{2n+3}} + \frac{1}{(\beta - \gamma)^{2n+3}} \right) + O \left( \frac{UM^{2n+2}}{T^{n+1}N^{2n+1} + M} \right) \right).
\]

Then by (3.2), the error terms in (3.1) are

\[
\int_\alpha^u g(x)e(f(x))dx = e(f(\gamma)) \int_{-r}^r e(\lambda_2 y^2)g(x) \frac{dx}{dy} dy = e(f(\gamma)) \sum_{k=0}^{2n} \varpi_k \int_{-r}^r e(\lambda_2 y^2) dy + e(f(\gamma)) \int_{-r}^r e(\lambda_2 y^2)Q(y) dy,
\]

with \(Q(y)\) defined in (1.12), because \(\varpi_k, 0 \leq k \leq 2n,\) is independent of \(y\). We will now compute the first integral on the right hand side of (3.3).

Lemma 3.1. For even \(k = 2j\) define

\[
\phi^{(k)}_i(y) = (-1)^{t-1}(2j-1)\cdots(2j-(2t-3))y^{2j-(2t-1)}, \quad 1 \leq t \leq j+1,
\]

where the numerator equals 1 for \(t = 1\). Then

\[
\int_{-r}^r e(\lambda_2 y^2) dy = \left[ e(\lambda_2 y^2) \sum_{i=1}^j \phi^{(k)}_i(y) \right]_{-r}^r + \frac{(-1)^j(2j-1)}{(4\pi\lambda_2)^{t-1}} \int_{-r}^r e(\lambda_2 y^2) dy.
\]

Proof. By (3.4) and (3.5) we know

\[
\phi^{(k)}_i(y) = \frac{y^{k-1}}{4\pi\lambda_2}, \quad \phi^{(k)}_i(y) = \frac{(\phi^{(k)}_{i-1}(y))'}{4\pi\lambda_2 i y} \text{ for } 2 \leq t \leq n+1
\]

for \(2n \geq k\). Applying integration by parts \(j\) times we get

\[
\int_{-r}^r e(\lambda_2 y^2) dy = \left[ e(\lambda_2 y^2) \sum_{i=1}^j \phi^{(k)}_i(y) \right]_{-r}^r - \int_{-r}^r e(\lambda_2 y^2)(\phi^{(k)}_j)(y) dy.
\]
From (3.4) for \( t = j \) we get

\[
\phi^{(k)}_j(y) = (-1)^{j-1} \frac{(2j-1)!! y^{2j}}{(4\pi\lambda_2)^j}, \quad (\phi^{(k)}_j(y))^{\prime} = (-1)^{j-1} \frac{(2j-1)!!}{(4\pi\lambda_2)^j}.
\]

Substituting (3.9) into (3.8) we prove (3.6). \( \square \)

The integral on the right side of (3.6) can be expressed in terms of the probability integral (cf. Gradshteyn and Ryzhik [2] 8.251.1)

\[
\Phi(x) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{x} e^{-t^2} dt.
\]

In fact,

\[
\int_{-\infty}^{x} e^{\lambda y^2} dy = \frac{2}{\sqrt{2\pi\lambda}} \int_{0}^{x} e^{it^2} dt = \frac{\Phi(x)}{\sqrt{2\lambda}} \Phi\left( \frac{x}{\sqrt{2\lambda}} \right)
\]

for \( x = r\sqrt{2\pi\lambda} \), by [2] 8.256.1. An asymptotic expansion of (3.10) is given by [2] 8.254. Therefore

\[
\int_{-\infty}^{r} e^{\lambda y^2} dy = \frac{\Phi(x)}{\sqrt{2\lambda}} \Phi\left( \frac{x}{\sqrt{2\lambda}} \right) + \sum_{k=0}^{d} \frac{(-1)^k (2k-3)!!}{(4\pi\lambda_2)^{k+1} r^{2k+2}} + O\left( \frac{1}{\lambda_2^{d+1} r^{2d+1}} \right).
\]

By (3.11) we see that the error term in (3.11) is

\[
O\left( \frac{M^{2d+2}}{T^{d+1} r^{2d+1}} \right).
\]

Substituting (3.11) and (3.12) into (3.2), we get for \( k = 2j \)

\[
\int_{-\infty}^{r} e^{\lambda y^2} y^k dy = \frac{(-1)^j (2j-1)!!}{(4\pi\lambda_2)^j} \frac{\Phi(x)}{\sqrt{2\lambda}} \Phi\left( \frac{x}{\sqrt{2\lambda}} \right) + 2e^{\lambda y^2} \frac{(-1)^j (2j-1)!!}{(4\pi\lambda_2)^j} \frac{1}{4\pi\lambda_2 r} \left( 1 + \sum_{k=2}^{d} \frac{(2k-3)!!}{(4\pi\lambda_2)^{k-1} r^{2k-2}} \right)
\]

+ \left[ e^{\lambda y^2} \sum_{i=1}^{j} \phi^{(k)}_i(y) \right]_{-r}^r + O\left( \frac{M^{2d+2+j}}{T^{d+j+1} r^{2d+1}} \right).
\]

By (3.1) and (3.2) we see that the second term is exactly equal to \( \sum_{i=j+1}^{j+1} \phi^{(k)}_i(r) \). Therefore

\[
\int_{-\infty}^{r} e^{\lambda y^2} y^k dy = \frac{(-1)^j (2j-1)!!}{(4\pi\lambda_2)^j} \frac{\Phi(x)}{\sqrt{2\lambda}} \Phi\left( \frac{x}{\sqrt{2\lambda}} \right) + \left[ e^{\lambda y^2} \sum_{i=1}^{j+1} \phi^{(k)}_i(y) \right]_{-r}^r
\]

+ O\left( \frac{M^{2d+2+j+2}}{T^{d+j+1} r^{2d+1}} \right).

For \( j \leq n + 1 \), we take \( d = n + 1 - j \) in (3.13).

Now we need a bound for \( \varpi_k \):

\[
\varpi_k \ll \sum_{i=0}^{k} \frac{U}{N^l} \sum_{j=1}^{k-l} \frac{M^{2j}}{T^j} \sum_{i=0}^{k} \frac{M^l}{M^{k-l+j+2}} \ll \frac{U}{M^k} \sum_{i=0}^{k} \left( \frac{M}{N} \right)^i \ll U \left( \frac{1}{N} + \frac{1}{M} \right) \text{ for } 1 \leq k \leq 2n
\]

by (1.11) and (1.4)–(1.6). Then by (3.14) and (3.15) we have

\[
\sum_{k=0}^{2n} \varpi_k \int_{-\infty}^{r} e^{\lambda y^2} y^k dy = \frac{\Phi(x)}{\sqrt{2\lambda}} \Phi\left( \frac{x}{\sqrt{2\lambda}} \right) \left[ e^{\lambda y^2} \sum_{i=1}^{n} \varpi_{2j-i} \frac{(-1)^{j-i} (2j-1)!!}{(4\pi\lambda_2)^j} \right]_{-r}^r + \varpi \left[ e^{\lambda y^2} \sum_{i=1}^{n} \varpi_i \phi^{(k)}_i(y) \right]_{-r}^r
\]

+ O\left( \sum_{i=0}^{n} U \left( \frac{M}{N} + 1 \right)^{2j} \frac{M^{2n-2j+4}}{T^{n+2+j+3}} \right),
where we added in the terms for odd \( k \) which are zero anyway.

Let us turn to the second integral on the right hand side of (3.3).

**Lemma 3.2.** With \( Q(y) \) as in (1.12) define

\[
\psi_1(y) = \frac{Q(y)}{(4\pi i \lambda_2)y} \quad \text{and} \quad \psi_k(y) = -\frac{\psi_{k-1}(y)}{(4\pi i \lambda_2)y} \quad \text{for} \quad 2 \leq k \leq n+3.
\]

Then

\[
\int_{-r}^{r} Q(y)e(\lambda_2 y^{2})dy = \left[ e(\lambda_2 y^{2}) \sum_{k=1}^{n+1} \psi_k(y) \right]_{-r}^{r} + O\left( U \frac{M^{2n+2}}{N^{2n+1} + M} \right).
\]

**Proof.** From (1.12) and (2.30) we get

\[
Q(y) \ll_{n} U|y|^{2n+1}\left( \frac{1}{N^{2n+1}} + \frac{1}{M^{2n+1}} \right).
\]

Similarly by (1.12) and (2.34) we see that for \( 1 \leq t \leq n+2 \)

\[
Q^{(t)}(y) = \frac{d^t}{dy^t} \left( g(x) \frac{dx}{dy} \right) - \sum_{k=0}^{2n-t} \frac{(k + t)!}{k!} \omega_{k+t} y^{k}.
\]

Therefore by (2.34) for \( 1 \leq t \leq n+2 \)

\[
Q^{(t)}(y) \ll U|y|^{2n+1-t}\left( \frac{1}{N^{2n+1}} + \frac{1}{M^{2n+1}} \right).
\]

Next we choose a real number \( \delta \approx \lambda_2^{-1/2} \) such that \( r/\delta \) is a power of 2. The total variation of \( Q(y) \) on \([-\delta, \delta]\) is

\[
V(Q(y)) \ll \int_{-\delta}^{\delta} \frac{|dQ|}{dy} dy \ll U\delta|^{2n+1}\left( \frac{1}{N^{2n+1}} + \frac{1}{M^{2n+1}} \right).
\]

From (1.12), (3.18), (3.20), by the Second Derivative Test (see Lemma 5.1.3 of [3], p.88), we have

\[
\int_{-\delta}^{\delta} Q(y)e(\lambda_2 y^{2})dy \ll \max_{-\delta \leq y \leq \delta} \left| Q(y) \right| + V(Q(y)) \ll U\delta|^{2n+1}\left( \frac{1}{N^{2n+1}} + \frac{1}{M^{2n+1}} \right)
\]

\[
\ll \frac{U}{T_{n+1}} \frac{M^{2n+2}}{N^{2n+1} + M}.
\]

We split the range \( \delta \leq y \leq r, -r \leq y \leq -\delta \) into intervals of the form \( t \leq y \leq 2t, -2t \leq y \leq -t \). By integration by parts we have

\[
\int_{t}^{2t} Q(y)e(\lambda_2 y^{2})dy = \left[ e(\lambda_2 y^{2}) \sum_{i=1}^{n+1} \psi_i(y) \right]_{t}^{2t} - \int_{t}^{2t} e(\lambda_2 y^{2})\psi_{n+1}'(y)dy.
\]

From (3.19) we see that for \( 1 \leq k \leq n+3 \)

\[
\psi_k(y) = \sum_{i=0}^{k-1} c_{ki} \frac{Q^{(i)}(y)}{(4\pi i \lambda_2)^i y^{2k-1-i}} \ll \frac{U|y|^{2n+2-2k}}{\lambda_2^{k}} \left( \frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right)
\]

\[
\ll \frac{U^{2n+2-2k}}{\lambda_2^{k}} \left( \frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right),
\]

since \( t \leq \left| y \right| \leq 2t \). Hence

\[
\psi_{n+1}'(y) = -4\pi i \lambda_2 y \psi_{n+2}(y) \ll \frac{U}{\lambda_2^{n+1}} \left( \frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right),
\]
and the total variation of $\psi'_{n+1}(y)$ on $[t, 2t]$ is

$$V(\psi'_{n+1}(y)) = \int_t^{2t} |\psi''_{n+1}(y)| dy = \int_t^{2t} | - 4\pi i \lambda_2 \psi_{n+2}(y) + (4\pi i \lambda_2 y)^2 \psi_{n+3}(y) | dy.$$  

By the First Derivative Test (Lemma 5.1.2 of [3], p.88) we get

$$\int_t^{2t} e(\lambda_2 y^2) \psi'_{n}(y) dy \ll \frac{\max_{t \leq y \leq 2t} \{\psi'(y)\} + V(\psi'(y))}{\lambda_2 t} \ll \frac{U}{\lambda_2^{3+2t^2} \left( \frac{1}{N^{2n+1}} + \frac{1}{M^{2n+1}} \right)}.$$  

Note that $\lambda_2 \gg T/M^2$ and by (3.22) we get

$$\int_t^{2t} Q(y)e(\lambda_2 y^2)dy = \left[ e(\lambda_2 y^2) \sum_{k=1}^{n+1} \psi_k(y) \right]_t^{2t} + O\left( \frac{UM^2}{T^{n+2} t^2} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \right).$$  

Summing over ranges with $t = 2^k \delta, k = 0, 1, 2, \ldots$, we get

$$\int_{\delta}^{r} Q(y)e(\lambda_2 y^2)dy = \left[ e(\lambda_2 y^2) \sum_{k=1}^{n+1} \psi_k(y) \right]_{\delta}^{r} + O\left( \frac{UM^2}{T^{n+2} \delta^2} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \sum_{k \geq 1} \frac{1}{2^{2k}} \right) = \left[ e(\lambda_2 y^2) \sum_{k=1}^{n} \psi_k(y) \right]_{\delta}^{r} + O\left( \frac{U}{T^{n+1}} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \right).$$  

Similarly

$$\int_{-r}^{-\delta} Q(y)e(\lambda_2 y^2)dy = \left[ e(\lambda_2 y^2) \sum_{k=1}^{n} \psi_k(y) \right]_{-r}^{-\delta} + O\left( \frac{U}{T^{n+1}} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \right).$$  

Now by (3.24), (3.25), and (3.26) we see that

$$\int_{-r}^{r} Q(y)e(\lambda_2 y^2)dy = \left[ e(\lambda_2 y^2) \sum_{k=1}^{n} \psi_k(y) \right]_{-r}^{r} + O\left( \frac{U}{T^{n+1}} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \right).$$  

By (3.23) we have the following trivial estimates

$$\left[ e(\lambda_2 y^2) \sum_{k=1}^{n+1} \psi_k(y) \right]_{-r}^{r} \ll \sum_{k=1}^{n+1} \frac{U \delta^{2n+2-2k}}{T^{n+1}} \left( \frac{1}{M^{2n+1}} + \frac{1}{N^{2n+1}} \right) \ll \frac{U}{T^{n+1}} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right).$$  

Therefore by (3.20) and (3.27) we prove the lemma. □

Now by (3.13), (3.17), we have

$$\int_{u}^{v} g(x)e(f(x))dx = e(f(\gamma)) \frac{e(\frac{1}{8})}{\sqrt{f''(\gamma)}} \sum_{j=0}^{n} \alpha_{2j} (\frac{1}{4\pi i \lambda_2})^j$$  

$$+ e(f(\gamma)) \left[ e(\lambda_2 y^2) \sum_{k=0}^{2n} \alpha_{2k} \sum_{i=1}^{n+1} \phi_i^{(k)}(y) \right]_{-r}^{r} + e(f(\gamma)) \left[ e(\lambda_2 y^2) \sum_{k=1}^{n+1} \psi_k(y) \right]_{-r}^{r} + O\left( \sum_{j=0}^{n} \frac{U(M+1)}{T^{n+2} \phi_{2n-2j+3}} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \right).$$  

Therefore by (3.20) and (3.27) we prove the lemma.
By \(2.7\), the first \(O\)-term in \(3.28\) is
\[
(3.29) \quad \ll \frac{U}{T^{n+1}} \left( \frac{M^{2n+1}}{N^{2n}} + M \right) + \frac{U M^{2n+4}}{T^{n+2}} \left( \frac{1}{(\gamma - \alpha)^{2n+3}} + \frac{1}{(\beta - \gamma)^{2n+3}} \right)
\]
Let
\[
F(y) = f(\gamma) + \lambda_2 y^2 = f(x), \quad G(y) = g(x) \frac{dx}{dy},
\]
\[
\theta_i(y) = \frac{G(y)}{2\pi i F'(y)}, \quad \theta_i(y) = -\frac{\theta_{i-1}(y)}{2\pi i F'(y)} \text{ for } 2 \leq i \leq n + 1.
\]
Now we want to show the following two equalities,
\[
(3.30) \quad \left[ e(F(y)) \sum_{i=1}^{n+1} \theta_i(y) \right]_r^r = e(f(\gamma)) \left[ e(\lambda_2 y^2) \sum_{k=0}^{n} \varpi_k \sum_{i=1}^{n+1} \phi_i^{(k)}(y) \right]_r^r
\]
\[
+ e(f(\gamma)) \left[ e(\lambda_2 y^2) \sum_{k=1}^{n+1} \psi_k(y) \right]_r^r
\]
\[
(3.31) \quad \left[ e(F(y)) \sum_{i=1}^{n+1} \theta_i(y) \right]_r^r = \left[ e(f(x)) \sum_{i=1}^{n+1} H_i(x) \right]_u^n,
\]
where \(H_i(x)\) is defined in \(1.22\). By \(1.12\) we see that
\[
G(y) = \sum_{k=0}^{2n} \varpi_k y^k + Q(y), \quad \theta_1(y) = \sum_{k=0}^{2n} \varpi_k \frac{y^{k-1}}{4\pi i \lambda_2} + \frac{Q(y)}{4\pi i \lambda_2 y^{2}}.
\]
By \(3.7\) and \(3.16\) we see that for \(1 \leq i \leq n + 1\)
\[
\theta_i(y) = \sum_{k=0}^{2n} \varpi_k \phi_i^{(k)}(y) + \psi_i(y).
\]
Therefore \(3.30\) is true.

To prove \(3.31\) we see that
\[
F'(y) = f'(x) \frac{dx}{dy} \text{ and } G(y) = g(x) \frac{dx}{dy},
\]
and hence
\[
\theta_1(y) = \frac{G(y)}{2\pi i F'(y)} = \frac{g(x)}{2\pi i f'(x)} = H_1(x),
\]
\[
\theta_i(y) = -\frac{\theta_{i-1}(y)}{2\pi i F'(y)} = -\frac{H_{i-1}'(x)}{2\pi i f'(x)} = H_i(x) \text{ for } 2 \leq i \leq n + 1.
\]
Now by induction \(3.31\) follows from the last two formulas and the correspondence between \(y = r\) and \(x = v\), and between \(y = -r\) and \(x = u\).

Then we can conclude from \(3.28, 3.29, 3.30\) and \(3.31\) that
\[
(3.32) \quad \int_u^v g(x)e(f(x))dx = \frac{e(f(\gamma) + \frac{1}{2})}{\sqrt{f''(\gamma)}} \left( g(\gamma) + \sum_{j=1}^{n} \varpi_{2j} \frac{(-1)^j (2j - 1)!}{(4\pi i \lambda_2)^j} \right)
\]
\[
+ \left[ e(f(x)) \sum_{i=1}^{n+1} H_i(x) \right]_u^n + O \left( \frac{U}{T^{n+1}} \left( \frac{M^{2n+2}}{N^{2n+1}} + M \right) \right)
\]
\[
+ U \left( \frac{M}{N} + 1 \right) \frac{M^{2n+4}}{T^{n+2}} \left( \frac{1}{(\gamma - \alpha)^{2n+3}} + \frac{1}{(\beta - \gamma)^{2n+3}} \right).
\]
At last, we consider the whole integral \((1.1)\). Since 
\[ \int_{\alpha}^{\beta} = \int_{u}^{v} + \int_{u}^{v}, \]
then by (3.2) and (3.32), we can prove (1.8).

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References

1. V. Blomer, R. Khan and M. Young, Distribution of mass of holomorphic cusp forms, *Duke Math. J.*, 162(14) (2013), 2609-2644.
2. I.S. Gradshteyn and I.M. Ryzhik, *Table of Integrals, Series, and Products*, 6th edition, Academic Press, 2000, San Diego. Online Errata: [http://www.mathtable.com/errata/Gr6_errata.pdf](http://www.mathtable.com/errata/Gr6_errata.pdf)
3. M.N. Huxley, *Area, Lattice Points, and Exponential Sums*, London Math. Soc. Monographs 13, 1996.
4. M. Jutila and Y. Motohashi, Uniform bound for Hecke \(L\)-functions, *Acta Math.*, 195(1) (2005), 61-115.
5. M. McKee, Haiwei Sun and Yangbo Ye, Improved subconvexity bounds for \(GL(2) \times GL(3)\) and \(GL(3)\) \(L\)-functions, preprint.
6. N. Salazar and Yangbo Ye, Spectral square moments of a resonance sum for Maass forms, submitted to *Frontiers Math. China*.
7. T.H. Wolff, *Lectures on Harmonic Analysis*, edited by I. Laba and C. Shubin, University Lecture Series 29, Amer. Math. Soc., Providence, 2003.

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