PAINLEVÉ IV AND A THIRD-ORDER VIEWPOINT

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Abstract. We take a third-order approach to the fourth Painlevé equation and indicate the value of such an approach to other second-order ODEs in the Painlevé-Gambier list of 50.

Introduction

Among the full list of six, the fourth Painlevé equation PIV is the most complicated for which all solutions are meromorphic throughout the complex plane. PIV has the form

\[
\dddot{w} = F(z, w, \dot{w})
\]

where the right side is rational in all three variables: it fails to be polynomial in all variables only in having \(w\) in the denominator; for its precise form, see below. On the one hand, this means that the standard (local) existence-uniqueness theorem for second-order ODEs applies to PIV with initial data in which \(w(a) \neq 0\) and \(\dot{w}(a)\) are specified; on the other hand, such a standard theorem does not apply to PIV with initial data involving \(w(a) = 0\). Further differentiation produces a third-order ODE of the form

\[
\dddot{w} = G(z, w, \dot{w}, \ddot{w})
\]

where the right side is polynomial in all variables (and \(\dddot{w}\) is absent). A standard (local) existence-uniqueness theorem for third-order ODEs applies to this equation: locally, there exists a unique solution \(w\) for which \(w(a), \dot{w}(a)\) and \(\ddot{w}(a)\) take specified values. This circumstance has its consequences for PIV, some of which we address. One reason for the simplification that arises upon passage to third order is that derivatives enter PIV only in the combination \(\dddot{w} - \dot{w}^2/2w\). This precise combination appears in a further dozen of the 50 canonical forms that stem from the analysis of Painlevé and Gambier as listed by Ince in [1]; the third-order approach may be profitably considered there also, as we illustrate in a couple of cases.

Painlevé IV

The precise form of the fourth Painlevé equation in the literature varies as to the naming of its parameters; we find it convenient to adopt the form

\[
(\text{PIV}) \quad \dddot{w} = \frac{\dot{w}^2}{2w} + \frac{3}{2}w^3 + 4zw^2 + 2(z^2 - \alpha)w - \frac{\beta^2}{2w}.
\]

This is the form taken in [1] when the equation appears as XXXI in the list of 50 canonical forms; upon the extraction of PIV as fourth in the list of six Painlevé equations, the parameter \(\beta^2\) is relabelled as \(-2\beta\). The ratio \((\dot{w}^2 - \beta^2)/2w\) on the right side of PIV is to be understood as a limit when appropriate: thus, if \(w\) vanishes at \(a\) then its first derivative at \(a\) satisfies \(\dot{w}(a)^2 = \beta^2\) or \(\dot{w}(a) = \pm \beta\); this convenience is the reason behind our choice of form.

It is natural to view PIV as a complex equation: when we do so, the differentiability of a solution entails its analyticity where defined; analyticity at an (isolated) zero involves the Riemann continuation theorem. We may instead view PIV as a real equation: where defined, each solution is then plainly thrice-differentiable (and better) away from its zeros; at each zero, the limit understanding of the ratio on the right side of the equation renders the second derivative continuous, higher differentiability following by the mean value theorem.
Calculation of the third derivative is most conveniently effected after rearrangement, thus
\[ 2w\dddot{w} = \dot{w}^2 + 3w^4 + 8zw^3 + 4(z^2 - \alpha)w^2 - \beta^2. \]
Upon differentiation away from zeros, \( 2\dddot{w} \) cancels from each side to yield
\[ 2w\dddot{w} = 12w^3\dot{w} + 24zw^2\dot{w} + 8w^3 + 8(z^2 - \alpha)w\dddot{w} + 8zw^2 \]
whence
\[ \dddot{w} = (6w^2 + 12zw + 4(z^2 - \alpha))\dot{w} + 4(w + z)w. \]
Now let \( a \) be an isolated zero of \( w \): if \( z \to a \) then \( w(z) \to w(a) = 0 \) and \( \ddot{w}(z) \to \ddot{w}(a) = \pm \beta \) so that \( \dddot{w}(z) \to 4(a^2 - \alpha)\ddot{w}(a) \); it follows that \( w \) is thrice-differentiable at \( a \) with
\[ \dddot{w}(a) = 4(a^2 - \alpha)\ddot{w}(a). \]
In other words, \( w \) continues to satisfy the foregoing third-order equation at isolated zeros.

**Theorem 1.** If \( w \) is a solution to PIV then \( w \) satisfies the third-order ODE\( \text{(PIV')} \)
\[ \dddot{w} = (6w^2 + 12zw + 4(z^2 - \alpha))\dot{w} + 4(w + z)w. \]
**Proof.** The proof precedes the statement, in which we assume the zeros of the solution to be isolated. \( \square \)

As mentioned in the Introduction, the form of PIV' guarantees that it satisfies the standard (local) existence-uniqueness theorem appropriate to third-order ODEs: locally, there exists a unique solution \( w \) to PIV for which \( w(a) \), \( \dot{w}(a) \) and \( \ddot{w}(a) \) assume arbitrarily specified values. As also mentioned in the Introduction, the standard existence-uniqueness theorem appropriate to second-order ODEs is not applicable to PIV itself: in general, there is no solution to PIV with \( w(a) = 0 \) unless \( \dot{w}(a) = \pm \beta \); in case \( \beta = 0 \), if \( w \) satisfies \( w(a) = 0 \) then \( \ddot{w}(a) = 0 \) and the standard second-order theorem would force \( w \) to vanish identically near \( a \) (which it need not do). At the risk of repetition, it is of course the case that not all solutions to PIV' satisfy PIV: regarding those that do, if \( w(a) \neq 0 \) and \( \dot{w}(a) \) are specified then \( \ddot{w}(a) \) must be as determined from PIV, while if \( w(a) = 0 \) then \( \ddot{w}(a) = \pm \beta \) necessarily.

Our passage to the third-order equation PIV' has applications to the Painlevé equation PIV itself. We first consider the special case in which \( \beta \) is zero.

**Theorem 2.** Let \( w \) be a solution to PIV in case \( \beta = 0 \). If \( w \) has an isolated zero at \( a \) then \( \ddot{w}(a) \neq 0 \).
**Proof.** Note that \( w \) is a solution to PIV' by Theorem 1 note also that the vanishing of \( w(a) \) implies the vanishing of \( \dot{w}(a) \). If \( \ddot{w}(a) \) vanishes too, then the standard existence-uniqueness theorem for the third-order equation PIV' forces \( w \) to vanish near \( a \); consequently, the zero at \( a \) cannot be isolated. \( \square \)

Let us now take both parameters \( \alpha \) and \( \beta \) to be zero, and write PIV\(_0\) for the resulting version of the fourth Painlevé equation, thus:
\[ \dddot{w} = \frac{\dot{w}^2}{2w} + \frac{3}{2}w^3 + 4zw^2 + 2z^2w. \]
View this as a real equation, with real solutions. Let the real solution \( w \) to real PIV\(_0\) have \( a \) as an isolated zero. By Theorem 2 it follows that either \( \dddot{w}(a) > 0 \) or \( \dddot{w}(a) < 0 \): in the former case, \( \dddot{w}(t) \) passes from strictly negative to strictly positive as \( t \) increases through \( a \); in the latter case, \( \dddot{w}(t) \) passes from strictly positive to strictly negative as \( t \) increases through \( a \). Now restrict to an open interval \( I \ni a \) in which \( a \) is the unique zero of \( w \). The rule
\[ f(t) = \begin{cases} -\sqrt{w(t)} & \text{if } I \ni t \leq a \\ +\sqrt{w(t)} & \text{if } I \ni t \geq a \end{cases} \]
then defines on $I$ a square-root $f$ of $w$ that satisfies the second-order ODE
\[ 4f''' = f(3f^2 + 2t)(f^2 + 2t). \]
Observe that, unlike PIV itself, this differential equation for $f$ obeys the standard existence-uniqueness theorem for second-order ODEs. For more detail on this, see particularly Theorem 5 in [2]. Incidentally, notice that Theorem 2 of the present paper justifies the discussion that surrounds Theorem 8 (and thereby supports Theorem 5) of [2] without invoking the (difficult) meromorphicity of solutions to PIV.

Twelve more ODEs

Passage to a third-order ODE has benefits not only for the fourth Painlevé equation but also for other second-order ODEs among the 50 canonical forms that result from the Painlevé-Gambier classification as listed in [1]. The third-order approach succeeds in part because derivatives enter PIV in the precise combination
\[ \ddot{w} - \frac{w^2}{2w} = \frac{2w\dot{w} - w^2}{2w} \]
and
\[ (2w\ddot{w} - \dot{w}^2)^* = 2w\dddot{w}. \]
Now, a further twelve equations in the list of 50 feature derivatives in just this combination; accordingly, analogous passage to the associated third-order ODE may be contemplated in these cases, too. We have already taken such an approach in [3], where we analyzed the relationship between the (homogeneous) second Painlevé equation and equation XX from the list of 50. In the present section, we focus primarily on the practical utility of passage to third order, which we illustrate by a couple more cases.

The thirty-second equation in the list of 50 has the following form:

\[ \text{(XXXII)} \]
\[ \ddot{w} = \frac{\dot{w}^2}{2w} - \frac{1}{2w} = \frac{\dot{w}^2 - 1}{2w}. \]
It is recorded in [1] that the substitution $w = u^2$ engenders a first integral (though the recorded first integral contains a minor misprint). In fact
\[ 2u^2 + 2u\dot{u} = \ddot{w} = \frac{\dot{w}^2}{2w} - \frac{1}{2w} = \frac{(2u\dot{u})^2}{2u^2} - \frac{1}{2u^2} = 2u^2 - \frac{1}{2u^2} \]
so that
\[ 2\dddot{u} = -\frac{\dot{u}}{2u^3} \]
which leads immediately to the first integral
\[ \dot{u}^2 = K + \frac{1}{4u^2}. \]
It is certainly possible to integrate further, so as to determine $u$ and thereby determine $w$. However, it is instructive to pass directly from XXXII to the associated third-order equation. Thus: rearrange to obtain
\[ 2w\dddot{w} = \dot{w}^2 - 1 \]
and differentiate to deduce
\[ 2w\dddot{w} = 0 \]
wherefrom it is immediate that $w$ is (at most) a quadratic
\[ w = az^2 + bz + c; \]
all that remains is to filter this solution to the third-order equation through XXXII itself, which yields the constraint $b^2 - 4ac = 1$. Passage to third-order has furnished a surprisingly direct route to the solution of this particular second-order equation.
The ‘m = 2’ case of the seventeenth equation in the list of 50 has the simpler form

\[ \ddot{w} = \frac{\dot{w}^2}{2w}. \]

We may approach this equation by first rearranging it as \(2\ddot{w}/\dot{w} = \dot{w}/w\) and so deducing the first integral \(\dot{w}^2 = Kw\) with an elementary solution for \(w\); alternatively, we may rearrange the equation as \(2(\dot{w}/w)^\cdot = -(\dot{w}/w)^2\) and proceed accordingly. Instead, we may pass directly to the associated third-order equation: thus, rearrangement and differentiation again lead to

\[ 2w \dddot{w} = 0 \]

and thence to the quadratic \(w = az^2 + bz + c\); satisfaction of \(\text{XVII}\) requires that \(b^2 - 4ac = 0\) whence the quadratic is a perfect square. Again, passage to third order has proved to be practically efficient.

As a last illustration of the third-order approach, we take the twenty-ninth equation in the list of 50:

\[ \text{XXIX} \]

\[ \ddot{w} = \frac{\dot{w}^2}{2w} + \frac{3}{2} w^3. \]

In this case, the rearrangement \(2w\ddot{w} = \dot{w}^2 + 3w^4\) leads after cancellation to the third-order equation

\[ \dddot{w} = 6w^2 \dot{w} = (2w^3)^\cdot \]

from which we deduce that

\[ \dot{w} = 2w^3 + k; \]

multiplication by \(2\dot{w}\) throughout generates the first integral

\[ \dot{w}^2 = w^4 + Kw + L \]

and filtration through \text{XXIX} shows that the constant \(L\) is zero.

For the handling of some equations, an alternative identity may rival

\[ (2\dddot{w} - \dot{w}^2)^\cdot = 2w \dddot{w}. \]

Explicitly, notice that

\[ \left(\frac{\dot{w}^2}{w}\right)^\cdot = 2\dddot{w} - \frac{\dot{w}^3}{w^2} = 2\dddot{w} \left(\frac{\dot{w}}{w} - \frac{\dot{w}^2}{2w}\right) \]

or

\[ \left(\frac{\dot{w}^2}{w}\right)^\cdot = \frac{\dot{w}}{w^2} (2w\dddot{w} - \dot{w}^2). \]

For example, application of this identity to \text{XXIX} leads promptly to

\[ \left(\frac{\dot{w}^2}{w}\right)^\cdot = 3w^2 \dddot{w} = (w^3)^\cdot \]

and so to the first integral

\[ \dot{w}^2 = w^4 + Kw. \]

While on the topic of this rival identity, we shall demonstrate its usefulness in solving equation \text{XXXII} with which we began this section. Substitution of \text{XXXII} into the rival identity gives

\[ \left(\frac{\dot{w}^2}{w}\right)^\cdot = 2\frac{\dot{w}}{w} \left( -\frac{1}{2w} \right) = -\frac{\dddot{w}}{w^2} = \left(\frac{1}{w}\right)^\cdot \]

whence (with a judicious naming of the constant of integration)

\[ \dot{w}^2 = 1 + 4aw \]
and we have arrived at a first integral. Rather than press on and integrate one last time we step back and take a further derivative, following the same apparently perverse inclination that led us to pass from the fourth Painlevé equation to third order: thus

\[ 2\dot{w}\ddot{w} = 4a\dot{w} \]

so that shortly \( \ddot{w} = 2a \) and \( w = az^2 + bz + c \) (with \( b^2 - 4ac = 1 \) as before).

We remark that this same rival identity is also effective in handling XVII, XVIII and XIX, along with other equations in the list of 50, to the enjoyment of which we leave the reader.

REFERENCES

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[2] P.L. Robinson, *Painlevé IV: roots and zeros*, arXiv 1609.02494 (2016).

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