The Earth Model Column Collaboratory (EMC$^2$) v1.1: An Open-Source Ground-Based Lidar and Radar Instrument Simulator and Subcolumn Generator for Large-Scale Models

Israel Silber$^1$, Robert C. Jackson$^2$, Ann M. Fridlind$^3$, Andrew S. Ackerman$^3$, Scott Collis$^2$, Johannes Verlinde$^1$, and Jiachen Ding$^4$

$^1$Department of Meteorology and Atmospheric Science, Pennsylvania State University, University Park, PA, USA
$^2$Environmental Sciences Division, Argonne National Laboratory, Argonne, IL, USA
$^3$NASA Goddard Institute for Space Studies, New York, NY, USA
$^4$Department of Atmospheric Sciences, Texas A&M University, College Station, TX, USA

Correspondence: Israel Silber (ixs34@psu.edu)

Abstract.

Climate models are essential for our comprehensive understanding of Earth’s atmosphere and can provide critical insights on future changes decades ahead. Because of these critical roles, today’s climate models are continuously being developed and evaluated using constraining observations and measurements obtained by satellites, airborne, and ground-based instruments. Instrument simulators can provide a bridge between the measured or retrieved quantities and their sampling in models and field observations while considering instrument sensitivity limitations. Here we present the Earth Model Column Collaboratory (EMC$^2$), an open-source ground-based lidar and radar instrument simulator and subcolumn generator, specifically designed for large-scale models, in particular climate models, but also applicable to high-resolution model output. EMC$^2$ provides a flexible framework enabling direct comparison of model output with ground-based observations, including generation of subcolumns that may statistically represent finer model spatial resolutions. In addition, EMC$^2$ emulates ground-based (and air- or space-borne) measurements while remaining faithful to large-scale models’ physical assumptions implemented in their cloud or radiation schemes. The simulator uses either single particle or bulk particle size distribution lookup tables, depending on the selected scheme approach, to perform the forward calculations. To facilitate model evaluation, EMC$^2$ also includes three hydrometeor classification methods, namely, radar- and sounding-based cloud and precipitation detection and classification, lidar-based phase classification, and a Cloud Feedback Model Intercomparison Project Observational Simulator Package (COSP) lidar simulator emulator. The software is written in Python, is easy to use, and can be straightforwardly customized for different models, radars and lidars.

Following the description of the logic, functionality, features, and software structure of EMC$^2$, we present a case study of highly supercooled mixed-phase cloud based on measurements from the U.S. Department of Energy Atmospheric Radiation Measurement (ARM) West Antarctic Radiation Experiment (AWARE). We compare observations with the application of EMC$^2$ to outputs from four configurations of the NASA Goddard Institute for Space Studies (GISS) climate model (ModelE3) in single-column model (SCM) mode and from a large-eddy simulation (LES) model. We show that two of the four ModelE3
configurations can form and maintain highly supercooled precipitating cloud for several hours, consistent with observations and LES. While our focus is on one of these ModelE3 configurations, which performed slightly better in this case study, both of these configurations and the LES results post-processed with \( \text{EMC}^2 \) generally provide reasonable agreement with observed lidar and radar variables. As briefly demonstrated here, \( \text{EMC}^2 \) can provide a lightweight and flexible framework for comparing the results of both large-scale and high-resolution models directly with observations, with relatively little overhead and multiple options for achieving consistency with model microphysical or radiation scheme physics.

1 Introduction

The representation of cloud processes in large-scale models is continuously advancing, conceptually, and in the level of details and complexity implemented in the micro- and macro-physical schemes (e.g., Lin et al., 2019; Cesana et al., 2019). These improvements are reflected in the accuracy of the resulting model output (e.g., Klein et al., 2013; Lin et al., 2019; Myers et al., 2021; Wang et al., 2019), yet results still show large inter-model variability (e.g., Zelinka et al., 2020). This variability results from, among other sources, model weaknesses concerning their ability to predict atmospheric state variables and processes such as cloud geometrical and optical thicknesses (e.g., Cesana and Waliser, 2016; Klein et al., 2013), formation and transition of marine stratocumulus clouds (e.g., Rémillard and Tselioudis, 2015; Lin et al., 2014; Cesana et al., 2019), and the formation and maintenance of supercooled water (e.g., Cesana et al., 2012; Tan and Storelvmo, 2016; Silber et al., 2019b).

Meaningful model evaluation benefits from a direct ("apples-to-apples") comparison with observations (e.g., Bodas-Salcedo et al., 2014; Suzuki et al., 2015). For the evaluation of cloud representation, model output is often compared with active remote-sensing measurements from instrumentation such as lidars and radars, which provide information on the spatial structure of clouds and some direct indications about active microphysical processes. However, model evaluation is challenging because of observational detectability constraints (e.g., signal extinction), and lack of retrievals or large uncertainties in some microphysical and atmospheric state quantities by these instruments, for example, hydrometeor number concentration or water content. In addition, spatial resolution differences between a model’s output and an observing instrument’s measurement resolution present an additional difficulty.

To bridge the gap between large-scale models such as weather or climate models and active remote-sensing observations, instrument simulators with different purposes have been developed to estimate observed parameters using model output. For example, the Cloud-resolving model Radar SIMulator (CR-SIM; Oue et al., 2020) was developed to emulate zenith-pointing and scanning radar and lidar variables using high-resolution model output, with considerations of hydrometeor shape and the resulting scattering calculations (see also Mech et al., 2020). The Cloud Feedback Model Intercomparison Project Observational Simulator Package (COSP; Bodas-Salcedo et al., 2011; Swales et al., 2018), on the other hand, was developed to operate over large-scale model output targeting satellite data as observational constraints, although expansions for the emulation of ground-based radars and lidars have been developed (e.g., Zhang et al., 2018; Kuma et al., 2020). Among a variety of reasons such as the demanding computation associated with the emulation of satellite measurements and the ability to use and output detailed data to and from simulators, COSP is typically implemented on-line into models’ code to facilitate output.
To account for spatial resolution discrepancies, which are typically accentuated in the case of large-scale models due to their coarser resolution, some model evaluation studies and forward simulators emulate a higher spatial resolution by generating subcolumns (e.g., Bodas-Salcedo et al., 2008, 2011; Chepfer et al., 2008; Klein and Jakob, 1999; Lamer, 2019; Stephens et al., 2010; Swales et al., 2018; Webb et al., 2001). Statistics calculated using multiple generated subcolumns, which are nominally faithful to the processed model’s physics, can be directly compared with the associated observations, thereby mitigating spatial resolution biases and errors.

Here we present the Earth Model Column Collaboratory (EMC\(^2\)), an open-source ground-based lidar and radar simulator and subcolumn generator, which is designed to operate over large-scale model output while being faithful to the physics implemented in models’ microphysics or radiation schemes but can also be applied to high-resolution model output. EMC\(^2\) enables detailed evaluation of atmospheric thermodynamic profile and cloud properties extracted from local, regional, and global simulation outputs against long-term ground-based, air- or space-borne datasets. The software is written in Python, allowing quick installation and providing customizable operation (for scattering calculations, etc.). In section 2, we briefly illustrate the Python code and software followed by a detailed description of the subcolumn generator, the forward calculations, and classification routines currently implemented in EMC\(^2\). In section 3 we demonstrate the use of EMC\(^2\) by comparing observations of a case study of highly-supercooled drizzling cloud over West Antarctica (see Silber et al., 2019a) with application of EMC\(^2\) to outputs from a large-eddy simulation (LES) and the NASA Goddard Institute for Space Studies (GISS) ModelE3 climate model (see Cesana et al., 2019) in single-column model (SCM) mode. Finally, section 4 provides a summary of the code features and case study demonstration.

2 EMC\(^2\) Description

2.1 Software Description

EMC\(^2\) depicts a workflow for comparing forward calculated radar or lidar variables generated from large-scale model output with radar or lidar measurements. Fig. 1 shows a flowchart example of this workflow for using EMC\(^2\) to compare ModelE3’s output with high spectral resolution lidar (HSRL) measurements. The workflow starts with the Model class incorporated within the emc2.core module. The Model class contains model output field namelists and default hydrometeor parameters (Table 1). Using Python’s class inheritance, EMC\(^2\) allows the creation of a custom class specifying a given model’s namelists and parameters (a ModelE3 class in this example), which ensures that the model output can be standardized and used by the other modules in EMC\(^2\). Once loaded through the Model class internal methods, model output data are stored within the Model object using the xarray dataset module (Hoyer and Hamman, 2017).

The Model object is then input to the subcolumn generator (sect. 2.2). Note that subcolumn generator processing can be practically skipped by setting the number of subcolumns (N\(_s\)) to 1. The results of the subcolumn generator are stored in the xarray dataset contained within the Model object. Here, we introduce the Instrument class. Similar to the Model class, the Instrument class contains relevant information about the instrument being simulated (some of which is listed in Table 2)
Figure 1. Standard workflow of EMC$^2$ utilization for direct comparison between observations and model output. The flowchart exemplifies the use of EMC$^2$ to compare the NASA GISS ModelE3 climate model output with corresponding high spectral resolution lidar (HSRL) measurements. $N_s$ designates the specified number of subcolumns.

as well as the single-particle and bulk scattering calculation LUTs (see sect. 2.3 and 2.4). Currently, zenith-pointing instrument properties and scattering calculation LUTs are available for various lidars and radars operated by the Department of Energy (DOE) Atmospheric Radiation Measurement (ARM) climate research facility. That is, the elastic micropulse lidar (MPL) and HSRL, both operating at a wavelength of 532 nm (e.g., Flynn et al., 2007; Eloranta, 2005), the 910 nm CL31 ceilometer (Morris, 2016), the 1064 nm HSRL elastic channel (see Razenkov and Eloranta, 2018), the elastic channel of ARM’s Raman lidar operating at 355 nm (e.g., Newsom, 2009), the C-band scanning ARM precipitation radar (C-SAPR; starting in EMC$^2$ v. 1.2; Widener and Bharadwaj, 2012), the X-band scanning ARM cloud radar (XSACR; Widener et al., 2012b), the Ka-band ARM zenith radar (KAZR; Widener et al., 2012a), and the W-band ARM cloud radar (WACR and M-WACR; Widener and Johnson, 2006). Similar to the Model class, the Instrument class allows EMC$^2$ to be tailored to other radars and lidars.
Table 1. Hydrometeor class parameter values implemented in EMC²

| Density [kg/m³] | cl  | ci  | pl  | pi  | Notes / references                  |
|----------------|-----|-----|-----|-----|-------------------------------------|
|                | 1000| 500 | 1000| 250 | as in ModelE3                       |
| Lidar ratio    | 18  | 24  | 5.5 | 24  | Thorsen and Fu (2015, Table 3); Nott and Duck (2011) |
| Lidar linear depolarization ratio | 0.03 | 0.35 | 0.10 | 0.40 | Based on sect. 3’s case; see Silber et al. (2019a) |
| $a (b)$ in terminal velocity power law | 3e-7 (2) | 700 (1) | 841.997 (0.8) | 11.72 (0.41) | $V = aD^b$; $a$ units m^{1-b}s^{-1};$ cf. Morrison and Gettelman (2008, Table 2) |

deployed at different sites, which does not confine the analysis of measurements and model output to specific ARM instruments or sites, as long as the required parameters and suitable scattering LUTs are provided. Thus, the various parameters mentioned in the next subsections (e.g., all parameters shown in Tables 1 and 2) as well as the LUTs, can be easily specified and set to match configurations and assumptions implemented in different large-scale models, as well as complex scattering models more commonly implemented in cloud-resolving and LES models.

Following the subcolumn generator process, the Instrument and Model objects are then input to the lidar (sect. 2.3) or radar (sect. 2.4) simulators (lidar simulator in the case of fig. 1). The forward calculation results are stored in the same xarray dataset in the Model object. Simulated hydrometeor classification (sect. 2.5) can be performed following the completion of the forward calculations and stored in the xarray dataset. For comparison and visualization of these results, EMC² uses the Atmospheric Community Toolkit (ACT; Theisen et al., 2020). Thus, a SubcolumnDisplay object, inherited from ACT’s Display object contains the necessary methods for quick visualization of the simulated instrument variables. In addition, the SubcolumnDisplay object also contains several internal methods for generating curtain and profile plots of observational and simulated data stored in the Instrument or Model objects, allowing masking of simulated signals below instrument detectability, for example. The figures presented in the next section (sect. 3) show examples of EMC²’s visualization capabilities. Finally, since the data are in the xarray dataset format, EMC² also contains all of xarray’s analysis and visualization capabilities for these simulated datasets.

EMC² incorporates a suite of unit tests for each function using the pytest testing tool (https://pytext.readthedocs.io/) to inspect the integrity and functionality of the code. These unit tests are combined with continuous code integration using TravisCI integration service (https://travis-ci.com/), which runs the unit tests every time a developer submits a pull request on GitHub. If the unit test passes with the developer’s changes to the code, then the changes are approved to be a part of EMC². Documentation is also automatically generated from the metadata strings in each subroutine to ensure that each part of the code is well documented.
Table 2. Radar instruments and some of their characteristics currently implemented in EMC². The ARM SGP, ENA, NSA, AWR, and MOS site abbreviations denote the Southern Great Plains, Eastern North Atlantic, North Slope of Alaska, the AWARE campaign (at McMurdo Station, Antarctica), and the Multidisciplinary Drifting Observatory for the Study of Arctic Climate (MOSAiC) Expedition, respectively. The calculations of the minimum detectable signal ($Z_{e_{\min}}$) for KAZR are based on the analysis in Silber et al. (2018a). All $Z_{e_{\min}}$ values correspond to 2 $s$ integration time except for the AWR and MOS WACR, the values of which correspond to 0.2 $s$ integration time.

| Parameter and ARM site | XSACR | KAZR | WACR | Sources / references |
|------------------------|-------|------|------|----------------------|
| Frequency [GHz]        | 9.71  | 34.86| 95.04| Widener and Johnson (2006); Widener et al. (2012a, b) |
| Index of refraction for water ($|K_w|^2$) | 0.93  | 0.88 | 0.84 | Widener and Johnson (2006); Widener et al. (2012a, b) |
| $Z_{e_{\min}}$ at 1 km |       |      |      |                      |
| SGP                    | N/A   | -51.5 | -46.0 | 8 year KAZR dataset analysis; Widener and Mead (2004) |
| ENA                    | N/A   | -56.5 | N/A  | Analysis of 3.5 year KAZR2 dataset analysis |
| NSA                    | N/A   | -48.5 | N/A  | 7.5 year KAZR dataset analysis |
| AWR                    | -30.0 | -45.5 | -40.0| Falconi et al. (2018); 1 year KAZR dataset analysis; Burns et al. (2016) |
| MOS                    | -30.0 | -41.6 | -40.0| Falconi et al. (2018); 1 year KAZR dataset analysis; Burns et al. (2016) |

2.2 Allocation of Hydrometeors to Subcolumns

The following simulator description assumes large-scale model convective and/or stratiform cloud scheme outputs containing four hydrometeor classes: cloud water (cl), cloud ice (ci), rain (precipitating liquid; pl), and snow (precipitating ice; pi). While these four hydrometeor classes are widely used in various microphysics schemes, we note that EMC² can be generally adapted to cases in which fewer or additional classes are used. We note that all the acronyms, abbreviations, and symbols used throughout this study are listed in Appendix A.

The subcolumn allocation and forward calculations in EMC² can be performed using two main approaches; the first follows the assumptions and general logic implemented in large-scale model radiation schemes while the second follows the assumptions and logic in model microphysics schemes. In short, the radiation scheme approach largely treats hydrometeor fractions in a generalized manner and utilizes bulk (hydrometeor population) scattering lookup tables (LUTs) generated using specific size distribution assumptions. The microphysics scheme approach, on the other hand, includes full integration of single-particle scattering LUTs using hydrometeor particle size distributions prognosed by models with consideration of sub-grid hydrometeor class fraction variability assumptions.

We note that the detailed description below of the radiation and microphysics approaches is congruent with the current implementation of these approaches in the GISS ModelE3 climate model. However, the core of these approaches is similar in other climate and Earth system models (ESMs), and EMC² can be easily adapted to fit specific variations in a model assumptions (see sect. 2.1). For example, the microphysics approach currently operates only on stratiform microphysics scheme output using a two-moment bulk scheme (Gettelman and Morrison, 2015, hereafter MG2) that has been implemented in climate models such as the Community Earth System Model version 2 (CESM2) Community Atmosphere Model version 6 (CAM6) (Danabasoglu et al., 2020), the Energy Exascale Earth System Model (E3SM; Golaz et al., 2019), and ModelE3.
Prior to the radar and lidar forward calculations, EMC\(^2\) generates subcolumns for each model output column. These subcolumns emulate a higher model spatial resolution, which partially reconciles the locality of ground-based measurements and allows a more robust statistical model evaluation. Subcolumns are generated and populated with hydrometeors from the top-down using the maximum-random overlap approach (Tian and Curry, 1989; Fan et al., 2011; Hillman et al., 2018) similar to the description by Lamer (2019, ch. 6). EMC\(^2\) translates hydrometeor fractions in the model grid to a binary set of hydrometeor-containing and hydrometeor-free subcolumn bins. That is, given a specified number of subcolumns \(N_s\) determined by the user, the total number of hydrometeor-filled subcolumn bins at model level \(h\) and time step \(t\) is equal to the rounded value of \(N_s \times f_{hyd}(h,t)\), where \(f_{hyd}\) is the volume fraction of a hydrometeor class (e.g., \(f_{cl}, f_{pi}\)) or a generalized hydrometeor fraction \(f_{gen}\) used in the model radiation scheme, at the same model level and time step. Here and henceforth, we assume for simplicity an SCM output (no horizontal coordinate dimensions), even though EMC\(^2\) can generally operate not only on SCM simulation output but also on global simulation output. Note that EMC\(^2\) does not calculate \(f_{gen}\), but can use if it is a model output field. The \(f_{gen}\) parameter is a generalization of cloud and precipitation fractions denoting air volume in which scattering and absorption by hydrometeors influence radiative calculations for a model. (In the case of ModelE3, for example, \(f_{gen}\) corresponds to cloud fraction where the layer-average cloud opacity exceeds that of precipitation, or vice versa.).

The following steps are applied in order to populate subcolumns (number 1, 2, ..., \(N_s\)) with hydrometeors:

1. Convective cloud hydrometeors (cl and ci) are allocated to the first subcolumns (lowest index; if \(f_{hyd}(h,t) > 0\)), thereby generating cloud-containing subcolumns with maximum convective geometrical cloud depths.

2. Stratiform cloud hydrometeors are allocated to subcolumn bins unoccupied by convective cloud hydrometeors following the maximum-random overlap approach. Thus, stratiform clouds at model level \(h\) are first randomly allocated to subcolumn bins with overlying stratiform clouds (at level \(h+1\)), followed, if necessary, by random allocation to subcolumns with cloud-free bins directly above at level \(h+1\). This order of processing where clouds are preferentially extended vertically conforms with assumptions that are often implemented in large-scale model radiative transfer calculations. (Note that our objective is to be faithful to model physics, even though in the case of the maximum-random overlap approach, it could produce, in certain cases, larger model biases compared to observations (e.g., Hillman et al., 2018)). During this hydrometeor allocation step, maximum overlap between liquid and ice cloud hydrometeors is maintained in grid cells containing both cloud phases. For example, if \(f_{ci}(h,t) > f_{cl}(h,t) > 0\), all cl-containing subcolumn bins belonging to the same grid cell will necessarily contain ci as well (mixed-phase), but there will be additional subcolumn bins in that grid cell containing only ci hydrometeors. In the case where overlying stratiform hydrometeors exist, the overlying layer phase is not a factor of consideration, such that a subcolumn bin containing ci may be located right above a subcolumn bin containing cl or both cl and ci (mixed-phase), and vice versa.

3. Convective and stratiform precipitating hydrometeors (pl and pi) are allocated to subcolumns without convective-stratiform restrictions, such that convective and stratiform precipitation may co-exist in the same subcolumn bin. Similar to the stratiform cloud allocation, precipitation is allocated with maximum-random overlap, i.e., precipitation is first extended vertically. Thus, either stratiform or convective cloud hydrometeors may exist in the same subcolumn bins as
convective and/or stratiform precipitating hydrometeors (for example, describing a scenario of snow generated by a stratiform cloud seeding an underlying drizzling stratocumulus cloud field). If after the vertical extension of overlying precipitation some subcolumn bins are still to be populated with precipitating hydrometeors (e.g., in a case where \( f_{pl}(h,t) = f_{pl}(h+1,t) = 0 \) and \( f_{pi}(h,t) > f_{pi}(h+1,t) > 0 \)), these hydrometeors are randomly allocated to cloudy grid cells of the same type (convective or stratiform), followed by random allocation of hydrometeors to cloud-free subcolumn bins. Similar to the allocation of stratiform cloud hydrometeors, maximum overlap is maintained between liquid and ice precipitation, and overlying precipitation phase is not a factor of consideration during the step of vertical extension of precipitation.

Once the subcolumns are populated with hydrometeors, per hydrometeor class except for stratiform cl in the case of the microphysics approach, hydrometeor mixing ratio is set in every hydrometeor-containing subcolumn bin by

\[
q_{hyd}(s,h,t) = \frac{\bar{q}_{hyd}(h,t)}{f_{hyd}(h,t)},
\]

where \( q_{hyd} \) and \( \bar{q}_{hyd} \) designate the mixing ratio of a hydrometeor class (e.g., \( q_{cl}, q_{ci} \)) in subcolumn bin \( s \) and a corresponding model grid cell mean, respectively. In the case of cl when using the microphysics approach, at every model level \( h \) and time step \( t \), \( q_{cl}(s,h,t) \) is randomly set in cl-containing subcolumn bins such that it would comply with the sub-grid variability gamma distribution described by Morrison and Gettelman (2008, eq. 8) while adjusting the values in the last cl-containing subcolumn bin such that hydrometeor mass is conserved (as in the case of other hydrometeor classes), i.e.,

\[
\sum_{i=1}^{N_s} q_{hyd}(i,h,t) \times N_{hyd}(s,h,t) = \frac{\bar{q}_{hyd}(h,t)}{f_{hyd}(h,t)}.
\]

We note that sub-grid scale variability of cloud water in ModelE3 is tied to the sub-grid scale variability of moisture rather than set at a fixed value as in Morrison and Gettelman (2008).

In stratiform hydrometeor-containing subcolumn bins, hydrometeor number concentration is set for every hydrometeor class by

\[
N_{hyd}(s,h,t) = \frac{\bar{N}_{hyd}(h,t)}{f_{hyd}(h,t)},
\]

where \( N_{hyd} \) and \( \bar{N}_{hyd} \) designate the number concentration of a hydrometeor class in subcolumn bin \( s \) and a corresponding model grid cell, respectively. EMC2 assumes that convective schemes do not diagnose \( \bar{N}_{hyd} \), and hence, this information is currently not produced by the simulator.

2.3 Forward Calculation of Lidar Variables

2.3.1 Microphysics Approach

In the microphysics approach, applicable only to stratiform hydrometeors, per hydrometeor diameter \( D \), EMC2 calculates the hydrometeor size distribution, \( \phi_{hyd}(D,s,h,t) \), defined by \( q_{hyd} \) and \( N_{hyd} \), fully consistent with the MG2 scheme (see Morrison et al., 2009, eq. 1-3). Using LUTs containing full Mie calculations for spheres (following Bohren and Huffman, 1983, Appendix A) of single particle extinction and backscatter efficiencies at lidar operating wavelength \( \lambda_l \) \( (Q_{e_{hyd}}(D,\lambda_l) \) and \( Q_{bs_{hyd}}(D,\lambda_l) \), respectively), the lidar particulate extinction cross-section \( (\alpha_{p_{hyd}}(s,h,t)) \) and backscatter cross-section \( (\beta_{p_{hyd}}(s,h,t)) \) are
calculated in every hydrometeor-bearing subcolumn bin by

\[
\alpha_{\text{phyd}}(s,h,t) = \frac{\pi}{4} \int_{D_{\text{min}}}^{D_{\text{max}}} \phi_{\text{hyd}}(D,s,h,t) Q_{\text{chyd}}(D,\lambda_l) D^2 dD \approx \\
\frac{\pi}{8} \sum_{i=1}^{N_D-1} \left( \phi_{\text{hyd}}(D_i,s,h,t) Q_{\text{chyd}}(D_i,\lambda_l) D_i^2 + \phi_{\text{hyd}}(D_{i+1},s,h,t) Q_{\text{chyd}}(D_{i+1},\lambda_l) D_{i+1}^2 \right) \Delta D_i,i+1
\]

\[
\beta_{\text{phyd}}(s,h,t) = \frac{\pi}{4} \int_{D_{\text{min}}}^{D_{\text{max}}} \phi_{\text{bs}}(D,s,h,t) Q_{\text{bshyd}}(D,\lambda_l) D^2 dD \approx \\
\frac{\pi}{8} \sum_{i=1}^{N_D-1} \left( \phi_{\text{bs}}(D_i,s,h,t) Q_{\text{bshyd}}(D_i,\lambda_l) D_i^2 + \phi_{\text{bs}}(D_{i+1},s,h,t) Q_{\text{bshyd}}(D_{i+1},\lambda_l) D_{i+1}^2 \right) \Delta D_i,i+1
\]

where we use the trapezoidal rule for discrete integration over a series of \(D\) values, which can be unevenly spaced by \(\Delta D_{i,i+1} = D_{i+1} - D_i\), while noting that \(D_1 = D_{\text{min}}\) and \(D_{N_D} = D_{\text{max}}\), where \(N_D\) is the number of diameters for which \(Q_{\text{chyd}}, Q_{\text{bshyd}}\), and \(\phi_{\text{hyd}}\) are calculated. In the case of the Mie calculations currently available in EMCM\(^2\), \(D_1 = 0.1\ \mu\text{m}, D_{N_D} = 1\ \text{cm}, and \(\Delta D_{i,i+1}\) is constant and equals 0.1 \(\mu\text{m}\). The complex refractive indices \(m_{\text{hyd}}(\lambda_l)\) used for liquid hydrometeors in the Mie calculations can be taken from Segelstein (1981, Table 1) or Rowe et al. (2020, for a temperature of -10 \(^\circ\text{C}\)). Refractive indices for ice hydrometers are taken from Warren and Brandt (2008). The Maxwell-Garnett equation (Bohren and Battan, 1980, eq. 1) for a mixture of ice and air is used to calculate the effective \(m_{\text{hyd}}\) for ci and pi based on the ice densities implemented in EMCM\(^2\) for ModelE3 (Table 1) relative to bulk ice density of 917 \(\text{kg/m}^3\).

The total \(\alpha_p\) and \(\beta_p\) \((\alpha_{\text{ptot}}(s,h,t)\) and \(\beta_{\text{ptot}}(s,h,t)\), respectively\) are calculated as the sum of each of these variables for cl, ci, pl, and pi. The lidar linear depolarization ratio (LDR) is estimated by weighting fixed LDR values (per hydrometeor class; see Table 1) with the relative contribution of \(\beta_{\text{phyd}}\) to \(\beta_{\text{ptot}}\).

The cumulative optical thickness (from the surface upward) at the base of a given subcolumn bin, \(\tau_{\text{hyd}}\), is calculated by

\[
\tau_{\text{hyd}}(s,h,t) = \sum_{i=2}^{h} \alpha_{\text{phyd}}(s,i-1,t) \Delta z(i-1,t),
\]

where \(\Delta z(h,t)\) denotes the geometrical thickness of model level \(h\) at time step \(t\) and \(\tau_{\text{hyd}}(s,h=1,t) = 0\). The total integrated optical thickness, \(\tau_{\text{tot}}(s,h,t)\), being the sum of \(\tau_{\text{hyd}}\) for cl, ci, pl, and pi, is used to estimate the level of full lidar signal attenuation (received signal not detectable by the simulated instrument), the value of which can be used to constrain comparisons between model output forward calculations and observations. Lidar signal extinction at visible wavelengths typically occurs at an optical thickness of 3-5 (e.g., Sokolowsky et al., 2020, fig. 4), and hence, EMCM\(^2\) assumes by default that the lidar signal is extinct at a level where \(\tau_{\text{tot}} = 4\) to produce a lidar signal extinction mask. We note that EMCM\(^2\) allows calculating \(\tau_{\text{hyd}}\) from the top-down (i.e., at the top of a given model layer), thereby enabling simulation of airborne and spaceborne lidar measurements.

In some cases, the observational dataset only consists of measurements made by elastic lidars that do not allow direct retrieval of \(\alpha_p\) and \(\beta_p\). In these cases, observations can be compared with a diagnostic attenuated \(\beta_{\text{ptot}}, \beta_{\text{ptot.att}}\), which is often referred
to as the normalized relative backscatter (NRB; Campbell et al., 2002) after several lidar signal corrections are applied. $\beta_{p_{tot,att}}$ is calculated here by

$$
\beta_{p_{tot,att}}(s, h, t) = [\beta_p(s, h, t) + \beta_m(s, h, t)]T_m^2(s, h, t)\exp^{-2\eta_{tot}(s, h, t)} \tag{4}
$$

where $\beta_m$ and $T_m^2$ are the molecular backscatter cross-section and the two-way transmittance calculated following Penndorf (1957), and $\eta$ is the multiple scattering coefficient. The value of $\eta$ is set to 1 by default, effectively implying no multiple scattering by hydrometeors or perfect multiple scattering treatment. Such an assumption is not realistic in the vast majority of cases even though effective $\eta$ values closer to 1 are more likely given the common narrow field-of-view of ground-based (and airborne) lidars. While a values of 1 is most likely unrealistic, it precludes the introduction of additional confounding factors stemming from the dependence of multiple scattering effects on the hydrometeor particle size distributions (e.g., Eloranta, 1998), which could significantly impact model-observation comparisons of $\beta_{p_{tot,att}}$ (not shown). Moreover, this dependence of multiple scattering on particle size distributions suggests that a fixed $\eta$ value smaller than 1 would impact the faithfulness of the lidar simulator to the model physics. That said, we note that $\eta$ can be manually set to other fixed values based on the physical assumptions made or certain empirical results (e.g., Winker, 2003), and that the determination of the extinction level based on $\tau_{tot}$ is independent of $\eta$.

### 2.3.2 Radiation Approach

With the radiation approach, applicable to both stratiform and convective cloud scheme output, forward calculations rely on bulk scattering LUTs. Therefore, this approach is more than two orders of magnitude faster than the microphysics approach, thereby rendering EMC² more suitable for the analysis of large model output datasets. Using this approach, a geometric cross-sectional area for each hydrometeor-bearing subcolumn bin is first calculated assuming geometric scatterers:

$$
A_{hyd}(s, h, t) = \frac{3}{4} q_{hyd}(s, h, t)\rho_a(h, t)\rho_b r_{e_{hyd}}(h, t), \tag{5}
$$

where $\rho_a$ is the density of air, $\rho_b$ is the bulk density of the hydrometeor class phase (1000 and 917 kg/m$^3$ is the case of liquid and solid water, respectively), and $r_{e_{hyd}}$ is the effective radius of a hydrometeor class in the model grid cell, provided as a model output field. The $\alpha_{p_{hyd}}$ and $\beta_{p_{hyd}}$ are then calculated by

$$
\alpha_{p_{hyd}}(s, h, t) = Q_{e, vol_{hyd}}(r^*_{e_{hyd}}(h, t), \lambda_l) A_{hyd}(s, h, t) \tag{6a}
$$

$$
\beta_{p_{hyd}}(s, h, t) = Q_{bs, vol_{hyd}}(r^*_{e_{hyd}}(h, t), \lambda_l) A_{hyd}(s, h, t), \tag{6b}
$$

where $Q_{e, vol_{hyd}}(\lambda_l)$ and $Q_{bs, vol_{hyd}}(\lambda_l)$ represent in this case bulk efficiencies per unit volume taken from LUTs, in which they are provided as function of $r_{e_{hyd}}$. In Eq. 6a and 6b, however, $Q_{e, vol_{hyd}}$ and $Q_{bs, vol_{hyd}}$ are functions of the adjusted effective
radius, \( r_{ehyd}^* \), which equals \( r_{ehyd} \) in all hydrometeor classes and cloud types except for stratiform cloud ice and snow. In these two cases, \( r_{ehyd}^* = r_{ehyd} \Phi_{hyd} \) and \( \Phi_{hyd} = \Phi_{hyd} \frac{\rho_{hyd}}{\rho_b} + (1 - \Phi_{hyd})(\frac{\rho_{hyd}}{\rho_b})^{\frac{1}{3}} \), where \( \rho_{hyd} \) is the hydrometeor class density (see Table 1), and \( \Phi_{hyd} \) is a constant fluffiness factor. \( \Phi_{hyd} \) is used such that a value of 0 represents an equivalent mass bulk sphere while a value of 1 represents a fluffy sphere with an equivalent maximum dimension as in Gettelman and Morrison (2015). In the case of ModelE3, for example, \( \Phi_{hyd} \) is set by default to an intermediate value of 0.5, and generally serves as one of many tuning parameters.

The default \( Q_{e,vol_{hyd}} \) and \( Q_{bs,vol_{hyd}} \) in eq. 6a and 6b (respectively) implemented in EMC\(^2\) were calculated using single particle full Mie calculations in the case of liquid hydrometeors, and single-particle scattering LUTs for a severely roughened 8-column ice aggregate (Yang et al., 2013) in the case of solid hydrometeors. These ice aggregate scattering calculations have been shown by Holz et al. (2016) to enable a closure between infrared Moderate-Resolution Imaging Spectroradiometer (MODIS; Platnick et al., 2003) and visible Cloud-Aerosol Lidar with Orthogonal Polarization (CALIOP; Winker et al., 2009) satellite ice optical thickness retrievals, and were included in the MODIS collection 6 (C6) cloud product (Platnick et al., 2017). In order to calculate the \( Q_{e,vol_{hyd}} \) and \( Q_{bs,vol_{hyd}} \) LUTs, we assumed the same gamma distribution parameters as those implemented in the C6 dataset (see Hansen, 1971, eq. 1), consistent with the bulk LUTs utilized by ModelE3’s radiation scheme.

Following the calculations of \( \alpha_{p_{hyd}} \) and \( \beta_{p_{hyd}} \), the total variables \( \alpha_{p_{tot}} \) and \( \beta_{p_{tot}} \) as well as \( \tau_{hyd} \), \( \tau_{tot} \), and \( \beta_{p_{tot,att}} \) are calculated as in the microphysics approach (sect. 2.3.1).

### 2.4 Forward Calculation of Radar Variables

#### 2.4.1 Microphysics Approach

When the microphysics approach is selected in EMC\(^2\), the first three radar moments are calculated for each hydrometeor class in every hydrometeor-bearing subcolumn bin; that is, the equivalent reflectivity factor \( (Z_{ehyd}) \), the mean Doppler velocity \( (V_{D_{hyd}}) \) and the Doppler spectral width \( (\sigma_{D_{hyd}}) \) as well as total radar moment variables \( (Z_{e_{tot}}, V_{D_{tot}}, \text{and } \sigma_{D_{tot}}) \). Full Mie calculation LUTs for the emulated radars (Table 2) are first used to calculate \( \beta_{p_{hyd}} \) at the radar operating wavelength \( \lambda_r \) following eq. 2b. The \( m_{hyd}(\lambda_r) \) used for liquid in the Mie calculations can be taken from Segelstein (1981, Table 1) or Turner et al. (2016, using a temperature of -10 °C), while \( m_{hyd}(\lambda_r) \) values for ice are taken from Mätzler (2006, ch. 5.3; using temperature of -10 °C).

\( Z_{ehyd} \) in every hydrometeor-containing subcolumn bin is then calculated (in linear units) using (Doviak and Zrnić, 1993, eq. 4.33)

\[
Z_{ehyd}(s,h,t) = \frac{\beta_{p_{hyd}}(s,h,t)\lambda_r^2}{\pi^3 |K_w|^2},
\]

where \( |K_w|^2 \) is the dielectric factor for water used in the raw radar observational processing (see Table 2). Using the resultant \( Z_{ehyd}, V_{D_{hyd}} \) is then calculated by implementing the hydrometeor class terminal velocities parametrization used in the MG2 scheme (cf. Morrison and Gettelman, 2008, Table 2). In the calculation of \( V_{D_{hyd}} \) we neglect the model grid cell vertical wind, \( w \), which predominantly has little impact on the \( V_{D_{hyd}} \) value, especially at coarser spatial and temporal resolutions typical
to large-scale models. Finally, $\sigma_{D_{\text{hyd}}}$ is calculated, while we note that beamwidth and turbulent broadening (e.g., Chen et al., 2018) are omitted from this calculation, but will be added in future work.

To allow a valid comparison between the forward calculations and observations, signal attenuation is considered in the calculation of the attenuated $Z_{e_{\text{tot,att}}}$:

$$Z_{e_{\text{tot,att}}}(s,h,t) = 10\log_{10}\left(Z_{e_{\text{tot}}}(s,h,t) - 2\left(Y_{\text{hyd}_{\text{tot}}}(s,h,t) + Y_{\text{gas}}(s,h,t)\right)\right).$$  \hspace{1cm} (8)

where $Y_{\text{gas}}$ and $Y_{\text{hyd}_{\text{tot}}}$ are the one-way integrated attenuation at the base of the subcolumn bin (in dB) due to atmospheric gases ($O_2$ and $H_2O$; see Ulaby et al., 1981, sect. 5.3-5.5) and all hydrometeors, respectively. $Y_{\text{hyd}_{\text{tot}}}$ is calculated using

$$Y_{\text{hyd}_{\text{tot}}}(s,h,t) = 10\log_{10}\left(\sum_{i=2}^{h} \alpha_{p_{\text{tot}}}(s,i-1,t)\Delta z(i-1,t)\right),$$  \hspace{1cm} (9)

where $\alpha_{p_{\text{tot}}}$ is determined by summing $\alpha_{p_{\text{hyd}}}$ based on eq. 2a calculated at $\lambda_r$ over all hydrometeor classes while setting $Y_{\text{hyd}_{\text{tot}}}(s,h=1,t) = 0$.

The vertical profile of the minimum detectable equivalent reflectivity factor, $Z_{e_{\text{min}}}$, is calculated by

$$Z_{e_{\text{min}}}(h,t) = Z_{e_{\text{min}}}(z=1000\text{ m}) + 20\log_{10}\left(z(h,t)/1000\right),$$  \hspace{1cm} (10)

where $z(h,t)$ is the height at the base of model level $h$ (in meters) at time step $t$ and $Z_{e_{\text{min}}}(z=1000\text{ m})$ is the minimum detectable signal at 1 km (using the highest sensitivity mode; Table 2). When compared with observations, subcolumn bins where $Z_{e_{\text{tot,att}}}(s,h,t) < Z_{e_{\text{min}}}(h,t)$ can be treated as returned signal below the radar noise floor, and hence, are effectively considered hydrometeor-free.

### 2.4.2 Radiation Approach

When the radiation approach is selected in EMC$^2$, forward radar calculations using bulk LUTs are limited to the zeroth radar moment ($Z_e$) due to a set of limitations:

1. Large-scale model radiation schemes are not informed with hydrometeor fall velocities. Moreover, fall velocity parametrizations in microphysics schemes do not necessarily fully overlap with the hydrometeor size and shape assumptions implemented in radiation schemes.

2. Noting that EMC$^2$ operates off-line, hydrometeor class fall velocities are typically reported in model outputs as weighted means. Because not all cloud schemes enable back-tracing of hydrometeor class fall velocities as a function of particle diameter using analytical expressions and weighted output fields (e.g., the convective cloud scheme in ModelE3; see Elsaesser et al., 2017), hydrometeor class fall velocities per subcolumn bin cannot be straightforwardly reproduced.

3. The total radar moments include combinations of the different hydrometeor class mixing ratios, and hence, cannot be determined using a single set of bulk LUTs per hydrometeor class.
Thus, EMC\textsuperscript{2} calculates only the $Z_{e\text{hyd}}$, $Z_{e\text{tot}}$, and $Z_{e\text{tot,att}}$. Eqs. 6b and 7 are used to calculate $Z_{e\text{hyd}}$ with bulk scattering LUTs for $\lambda_e$ based on full single particle Mie scattering calculations in the case of liquid hydrometeors and single particle scattering calculations for the Yang et al. (2013) 8-column aggregate at a temperature of 270 K (see Ding et al., 2017) in the case of solid hydrometeors. Similar to the implementation of the radiation approach in the forward lidar calculations (sect. 2.3.2), $r_{e\text{hyd}}$ considers the fluffiness factor in the case of solid (ice) hydrometeors. Finally, $Z_{e\text{tot}}$, and $Z_{e\text{tot,att}}$ are calculated similar to the microphysics approach using eq. 8 and 9 with $\alpha_{p\text{hyd}}$ (and $\alpha_{p\text{tot}}$) calculated using eq. 6a.

2.5 Hydrometeor Classifications

Once the total lidar and/or radar variables are calculated, EMC\textsuperscript{2} can be used to classify the subcolumn simulator output. Classification masks can serve as tools for direct comparisons between the simulator output and observational data utilizing similar classification methodologies, some of which can be used to calculate water phase ratios. Currently, EMC\textsuperscript{2} incorporates three hydrometeor classification methods: the radar-sounding cloud and precipitation detection and classification, the modified fixed lidar variable threshold phase classification, and the COSP lidar simulator emulator (henceforth referred to as the COSP emulator; 'emulator' is used here in its generic sense rather than a machine-learning context, and there is no training involved).

The radar-sounding cloud and precipitation detection method (Silber et al., 2021a; see also Vassel et al., 2019) emulates the combined use of relative humidity with respect to water (RH) sounding measurements for the detection of liquid-bearing clouds and radar echoes (received signal above the radar noise floor) for the detection of precipitating hydrometeors. In the case of large-scale model output, RH below 100% in a model grid cell does not necessarily indicate a lack of cloud water, because of implemented assumptions concerning the sub-grid distribution of cloud water content (e.g., Smith, 1990). Therefore, the approach most consistent with the observational method is to simply use the cloud water-bearing subcolumn bins (see sect. 2.2) to classify the subcolumn bin as "cloud". "Precipitation" are those subcolumns bins in which $Z_{e\text{tot,att}} \geq Z_{e\text{min}}$. Subcolumn bins that can be classified as both "cloud" and "precipitation" are set as "mixed". We note that at temperatures below 0 °C, the "mixed" classification type becomes more likely to represent a mixed-phase cloud with decreasing temperatures, but in general, may represent bins containing only liquid hydrometeors.

The modified fixed lidar variable threshold phase classification method is similar to previous studies that incorporated fixed LDR and $\beta_{p\text{tot}}$ thresholds to classify hydrometeor-bearing air volumes to "liquid" and "ice-only" using lidar measurements (e.g., Shupe, 2007; Thorsen and Fu, 2015). By default, however, EMC\textsuperscript{2} includes two additional "undefined" classes that cover intermediate regions in the LDR-$\beta_{p\text{tot}}$, such that a subcolumn bin classified as "undefined1" has a higher probability that it includes some amount of liquid water while "undefined2" is more likely to only contain ice hydrometers (see sect. 3.3 for discussion and illustration of the default thresholds). The notion behind the addition of these two "undefined" classes is the fixed-threshold method limitations that could originate in:

1. Drizzle- or rain-bearing air-volumes, which may produce moderate $\beta_{p\text{tot}}$ and LDR on the order of 0.1, especially when ice hydrometeors are present in the same air volumes (e.g., Derr et al., 1976; Sassen, 2003; Silber et al., 2019a).
2. Cases in which an entire relatively tenuous liquid-bearing cloud layer or other cases with liquid-bearing air-volumes just above cloud base occur concurrently with ice precipitation sufficiently concentrated and intense to generate a combination of high $\beta_{p_{\text{tot}}}$ and LDR that can reach values of 0.10-0.20 (e.g., Derr et al., 1976; Silber et al., 2018b, fig. S1). This influence of ice hydrometeors also applies to rain-bearing air volumes.

3. Horizontally-oriented ice hydrometeors that may produce low (moderate) LDR ($\beta_{p_{\text{tot}}}$) via specular reflection even in cases when the lidar is titled up to several degrees off-zenith (e.g., Noel et al., 2002; Silber et al., 2018b, Appendix A). Note that this limitation only applies to observations, since specular reflection and ice particle canting angles are not represented in large-scale models.

Adaptive fixed thresholds that vary with site, instrument, and period of study (e.g., Silber et al., 2018a; Zaremba et al., 2020) or lidar ratio constraints (e.g., Thorsen and Fu, 2015) can compensate for some of these limitations. However, these approaches cannot be objectively translated to the model output domain to enable a direct comparison between the observations and the simulator output. Therefore, the modified fixed threshold routine, which largely agrees with existing measurements yet acknowledges both model and observational uncertainties may allow better direct comparisons to be made.

The emulator of the COSP lidar simulator follows the same equations and logic of the on-line lidar simulator (Cesana and Chepfer, 2013) implemented in numerous climate models. In short, the attenuated total backscatter (ATB) calculated in the COSP emulator routine while assuming $\eta = 0.7$ is used to calculate the lidar scattering ratio (the ratio of total to molecular attenuated backscatter) for the detection of hydrometeors in subcolumns by selecting scattering ratio values larger than 5. Calculated cross-polar ATB as a function of the total ATB is then used to classify the detected hydrometeors into liquid or ice, based on an empirical phase discrimination line. As the last step of this classification method, hydrometeors below (top-down lidar view) or above (bottom-up lidar view) a subcolumn bin with scattering ratio larger than 30 are classified as "undefined". Note that, unlike the on-line COSP lidar simulator, this emulator operates using the model vertical levels and does not interpolate the model output onto an evenly-spaced vertical grid. As with all other EMC$^2$ forward calculations and classification routines, this emulator can operate using a top-down viewing approach thus providing a bridge between the COSP simulator and EMC$^2$.

Finally, EMC$^2$ also includes internal functions to calculate mass or frequency phase ratios using each of these hydrometeor classification methods, providing metrics to compare model output with observations or with outputs from other models.

3 Case Study Example: Highly Supercooled Antarctic Cloud

To demonstrate the application of EMC$^2$ and its output using the different forward calculation approaches, here we describe and analyze a Lagrangian LES case study (Silber et al., 2019a) adjusted for running and testing the ModelE3 climate model (as well as other climate models) in SCM mode.
3.1 Case Description

As described by Silber et al. (2019a), the stratiform cloud event that we compare with model simulations was observed over McMurdo Station, Antarctica, as part of the ARM West Antarctic Radiation Experiment (AWARE; Lubin et al., 2020), on August 16, 2016. During the event, cyclone-driven wind confluence with southwesterly katabatic flow resulted in relatively warm and moist marine air convergence along the Ross Ice Shelf coast, part of which was advected towards the McMurdo Station measurement site (fig. 2a). This air convergence induced a widespread (>1000 km) cloud field evident by Clouds and the Earth’s Radiant Energy System (CERES) measurements indicating an extensive region with enhanced top of atmosphere (TOA) upwelling longwave radiation (dashed green shape in fig. 2b; note that a surface-based temperature inversion is common during the austral winter resulting in smaller TOA radiation fluxes).

Over McMurdo Station, a decoupled persistent mixed-phase cloud with temperatures as low as -29 °C was observed for ~39 hours. The observed cloud was nearly continuously precipitating ice particles and was also drizzling for more than 7 hours, concluded from a comprehensive analysis of sounding, HSRL, and KAZR measurements (see Silber et al., 2019a).

3.2 ModelE3 SCM Configuration

Based on Lagrangian simulations constrained by the remote-sensing observations, Silber et al. (2019a) postulated that the activated ice nucleating particle (INP) and cloud condensation nuclei (CCN) concentrations during the event were on the order of \(0.2 \, \text{L}^{-1}\) and \(20 \, \text{cm}^{-3}\), respectively, to enable drizzle to be produced and precipitate along with ice precipitation below the highly supercooled cloud base. Their simulations were 9 h in duration, starting on August 16, 2016, at 01:00 UTC and ending at 10:00 UTC. By imposing large-scale vertical wind extracted from back-trajectory calculations, they emulated the transport of the cloud layer, initially forming in a stable layer, towards McMurdo Station. The end of the simulation at 10:00 UTC designated the time at which the cloud field reached the fixed observational site at McMurdo, and hence, statistics of that hour of observations (10:00-11:00 UTC) were compared with the model output.

Here, we slightly adjusted the case study initialization files to enable running this case in a climate model’s SCM mode, while using the same single-hour of radar and lidar observations and the Distributed Hydrodynamic Aerosol and Radiative Modeling Application (DHARMA) model (Stevens et al., 2002) baseline LES output (see Silber et al., 2019a) as benchmarks. Namely, we simplified the profiles of vertical motion (as in Silber et al., 2020) and then converted the height coordinates of the initial sounding (see Silber et al., 2019a) and vertical wind time series to pressure coordinates (these converted sounding and forcing files are available at http://dx.doi.org/10.17632/gz4gdn3jvz.1). The utilization of these files to run the highly supercooled cloud case study enables testing of an SCM, and hence, the cloud schemes implemented in a climate model.

The LES is initialized here with activated INP concentration of \(0.1 \, \text{L}^{-1}\) and with both cloud ice and snow hydrometeor classes (whereas Silber et al. (2019a) used only a single ice class), and the simplified vertical motion profiles are then imposed. These few LES adjustments make the model configuration more consistent with climate model microphysics and SCM initialization while effectively resulting in the same hydrometeor content and cloud evolution as the baseline simulation presented.
Figure 2. (Top) Sea level pressure (contours; land areas are masked), total condensate water path (logarithmic color-scale), and surface winds (quivers) resolved by ERA5 (Hersbach et al., 2020) for August 16, 2016, at 10:00-11:00 UTC. The magenta-filled marker denotes McMurdo Station. The dashed green shape highlights a widespread cloud field along the Ross Ice Shelf (yellow shape) coast associated with the ground-based observations. The inset panel shows a topographic map of Antarctica (the red box highlights the region depicted in the main panel). (Bottom) Top of atmosphere upwelling longwave radiation measured by CERES Aqua on August 16, 2016, at 10:44-10:50 UTC. The 1 arc-minute topographic data were developed by the National Geophysical Data Center (NGDC Amante, 2009) and are freely available at https://www.ngdc.noaa.gov/mgg/global/global.html. Reproduced from Silber et al. (2019a).

in Silber et al. (2019a) (not shown). Thus, the SCM is run equivalently to the LES; that is, using the same initial sounding and forcing files, setting the coordinates to -77.85 °S, 166.72 °E, and prescribing a monomodal log-normal aerosol particle concentration of 20 $cm^{-3}$ with a mean radius of 0.1 $\mu m$, geometrical standard deviation of 2, and a hygroscopicity parameter of 0.4. We note that activated INP concentrations are not prescribed in the SCM simulations because (a) ModelE3’s final configurations are defined by specific values of certain model tuning parameters (among others) associated with INP param-
eterization, and (b) diagnostically prescribing the activated INP concentration is not faithful to the temperature-dependent approach implemented in the model, and hence, would not necessarily be informative of true climate model weaknesses.

3.3 Comparison Between Observations and ModelE3 SCM Using EMC$^2$

The following figures show the EMC$^2$ forward calculation results using the DHARMA LES simulation and ModelE3’s configuration Tun3, one of four configurations of ModelE3 derived in part via a machine learning parameter tuning approach that will be described in a manuscript in preparation, to be included in the Coupled Model Intercomparison Project Phase 6 (CMIP6). The SCM using this configuration was able to generate a cloud-top inversion and turbulent layer via cloud-top radiative cooling, and produced the best agreement with the observations and the LES relative to the other three configurations (see Appendix B). Whereas here we examine application of EMC$^2$ to a single ModelE3 configuration in SCM mode in a case where we can also compare with LES, we note that EMC$^2$ is designed to enable detailed evaluation of atmospheric thermodynamic profile and cloud properties extracted from global simulations of ModelE3 configurations and other climate models against long-term datasets at fixed sites in future dedicated work.

The left panels in Fig. 3 show the mixing ratios of the four hydrometeor classes evolving through the simulated SCM case study. These depicted mixing ratios are the output of the SCM’s stratiform cloud scheme, but because this simulation does not generate any convective hydrometeors (as expected), they also represent the total mixing ratios. Cloud water mass (top panel) dominates over the other hydrometeor classes in hydrometeor-bearing model grid cells through much of the simulation, even at lower levels in which the cloud water fraction is rather low (fig. 3, right panels). Rain (effectively drizzle) is produced by the model as well but has a relatively smaller mass compared with that of snow generated. This reduced amount of rainwater in the SCM simulation relative to rainwater dominance in the simulations of Silber et al. (2019a) is largely the result of the different autoconversion parameterization schemes implemented in ModelE3 (Seifert and Beheng, 2001) relative to that implemented by default in the DHARMA LES (Khairoutdinov and Kogan, 2000), which produces significantly smaller rain mass mixing ratios in this case (not shown), contrary to some previous studies (e.g., Heiblum et al., 2016; Xiao et al., 2021). Understanding the source of this differing autoconversion parameterization behavior requires a dedicated study that is beyond the scope of this manuscript.

Figs. 4 and 5 depict the HSRL and KAZR variables observed during a single hour over McMurdo Station and simulated with EMC$^2$ using the DHARMA LES three-dimensional output at the end of that simulation (without using the subcolumn generator), and by applying each of the three approaches on ModelE3 configuration Tun3 SCM output for 05:00 UTC. Because our goal in this section is to demonstrate that EMC$^2$ can reasonably match cloud observations given comparable input, we present the EMC$^2$-processed SCM output 4 hours into the simulation when cloud top heights are similar to observed instead of the end of the simulation (the SCM develops the supercooled cloud faster than the baseline LES; see Appendix B). When evaluating the processed model output against the observations, we essentially exchange temporal resolution with spatial resolution (three-dimensional model domain in the case of the LES) or an emulated spatial resolution (in the case of the
SCM). We set the number of SCM subcolumns \((N_s)\) to 100, which effectively represents hydrometeor class fractions up to the second decimal point, and enables drawing more robust statistics by emulating sub-grid variability of all hydrometeor classes combined.

The processed LES output exhibits good apparent agreement with the observations, evident by the comparable lidar and radar variable values and their horizontal variability (Figs. 4 and 5), the vertical cloud structure and boundaries, as well as the full lidar signal attenuation near cloud top (Fig. 4). A multi-layer cloud water structure developed by the LES is suggested by the intermittent breaks in the large \(\beta_{ptot}\) and \(\alpha_{ptot}\) values, supported by a separate analysis of \(\beta_{phyd}\) and \(\alpha_{phyd}\) (not shown). This multi-layer structure is also indicated by the lidar observations (Fig. 4), and was comprehensively discussed by Silber et al. (2019a).
Figure 4. Observed and simulated HSRL particulate backscatter cross-section ($\beta_{p,tot}$; top row), particulate extinction cross-section ($\alpha_{p,tot}$; middle row), and linear depolarization ratio (LDR; bottom row) using the DHARMA LES output and ModelE3 configuration Tun3 SCM output. The columns show (from left to right) the observations from McMurdo Station between 10:00-11:00 UTC, DHARMA LES three-dimensional output (horizontal dimensions stacked onto a single domain column dimension) processed using EMC$^2$ microphysics with the subcolumn generator turned off ($N_s = 1$), ModelE3 EMC$^2$ output using the microphysics approach, and EMC$^2$ output using the radiation approach. The DHARMA LES output corresponds to 10:00 UTC (arrival of cloud field at McMurdo Station; see sect. 3.2), while ModelE3’s EMC$^2$ panels depict the SCM output for 05:00 UTC (see text) processed with EMC$^2$ using 100 subcolumns. A mask denoting full lidar signal attenuation generated using a total accumulated optical thickness ($\tau_{tot}$) condition of $\tau_{tot} > 4$ is plotted over the simulated data (hatched areas).
Figure 5. As in fig. 4, but with the KAZR attenuated equivalent reflectivity factor ($Z_{\text{e tot,att}}$; top row), mean Doppler velocity ($V_{D\text{tot}}$; middle row), and spectrum width ($\sigma_{D\text{tot}}$; bottom row). Note that only $Z_{\text{e tot,att}}$ is calculated in EMC$^2$’s radiation approach (see sect. 2.4.2. A radar signal-to-noise ratio (SNR) mask of $Z_{\text{e min}} > Z_{\text{e tot}}$ ($Z_{\text{e min}}$ is the minimum detectable signal on each model level) is applied to the plotted data (hatched areas).

Using the microphysics approach, the SCM sub-grid variability is more pronounced relative to the radiation approach owing to the implementation of cloud water sub-grid variability (as defined in the MG2 microphysics scheme), which is not considered in the radiation approach. Evaluation of heights with full cloud cover, indicated by the large $\beta_{p\text{tot}}$ and $\alpha_{p\text{tot}}$ values (see also fig. 3, right), suggests that the SCM has an apparent reasonable agreement with the observations there. ModelE3’s macrophysics
scheme leads to notably ragged liquid cloud base heights that are evident in the SCM fields compared to the more uniform liquid cloud base in both observations and LES, as discussed further below. Relatively enhanced $\alpha_{p_{t_{\text{tot}}}}$ values are produced in some SCM subcolumns, which leads to full lidar signal attenuation, as also indicated by the HSRL measurements (fig. 4). Full lidar signal attenuation does not occur in this case using the radiation approach because of the uniform distribution of cloud water between subcolumns commensurate with $\tau_{t_{\text{tot}}}$ values just below the full attenuation threshold of 4 (not shown). The subcolumns with full lidar signal attenuation in the microphysics approach call for the use of radar measurements for cloud-top detection (e.g., fig. 5 discussed below).

Except for the overestimated LDR values below the supercooled cloud layer relative to the observations, the values of the other two lidar variables and the general scenario structure appear to agree with the observations and LES in both the microphysics and radiation approaches, which exhibit an encouraging consistency with each other (fig. 4). The comparison of the subcolumn-averaged lidar-variable profiles illustrated in fig. 6 allows a more quantitative comparison. This figure better indicates that the DHARMA LES $\alpha_{p_{t_{\text{tot}}}}$ within the main liquid-bearing cloud layer is within the range of the observations and their uncertainty, consistent with the conclusions of Silber et al. (2019a) that activated CCN number concentrations were low during this highly supercooled drizzle event.

The EMC$^2$ lidar output also highlights a few model weaknesses. For example, using the microphysics approach, cloud base (at a height of $\sim$1.0 km) is highly variable and may extend nearly down to the surface, which is in contradiction to the observations, where cloud base variability is on the order of a few hundred meters throughout the depicted period (fig. 4). Using the radiation approach, on the other hand, breaks in hydrometeor cover are seen below the fully overcast layer, as a result of the generalized hydrometeor fraction used in the radiative transfer calculations, which can be lower or higher than the associated hydrometeor class fraction (e.g., fig. 3, right). Lesser (greater) generalized hydrometeor fractions relative to the associated hydrometeor class fraction therefore implies greater (smaller) subcolumn bin mixing ratios (see eq. 1), and hence, enhanced (diminished) associated $\beta_{p_{\text{hyd}}}$ and $\alpha_{p_{\text{hyd}}}$ values. Thus, while the occurrence of low-level cloud water-bearing bins produces full attenuation of the simulated lidar signal in some subcolumns using the microphysics approach (fig. 4), the smaller $\alpha_{p_{\text{hyd}}}$ ensuing from the larger low-level generalized cloud fraction relative to the cloud water fraction (see fig. 3, right) causes no low-level (or any level in this case) signal extinction when the radiation approach is used (see fig. 4). Note that in both the microphysics and radiation approaches the subcolumn representation of hydrometeor mass remains consistent with the model output variables, i.e., eq. 1 holds for each hydrometeor class.

The microphysics and radiation approaches exhibit $Z_{e_{t_{\text{tot}},\text{att}}}$ values that are too large, especially at higher levels (figs. 5 and 6). EMC$^2$’s radar processing using the microphysics approach provides the $V_{D_{t_{\text{tot}}}}$ and $\sigma_{D_{t_{\text{tot}}}}$ variables in addition to the $Z_{e_{t_{\text{tot}},\text{att}}}$ variable. We note that because spectral broadening terms other than the microphysical broadening are currently not considered in EMC$^2$, the calculated $\sigma_{D_{t_{\text{tot}}}}$ generally serves as a lower bound for comparison purposes, i.e., the simulated $\sigma_{D_{t_{\text{tot}}}}$ values need to be smaller to some extent than the observed values. As indicated from fig. 5, both $V_{D_{t_{\text{tot}}}}$ and $\sigma_{D_{t_{\text{tot}}}}$ show grossly reasonable correspondence between the observations and the SCM, except at heights above $\sim$0.8 km, where $V_{D_{t_{\text{tot}}}}$ values show large deviations (fig. 6). A separate analysis (not shown) suggests that, in this case, these deviations are mainly the result of
**Figure 6.** Mean profiles of observed (time-averaged) and simulated (subcolumn-averaged) HSRL and KAZR variables calculated using EMC²’s microphysics (top) and radiation (bottom) approaches as shown in figs. 4 and 5 (full lidar signal attenuation and radar SNR masks are applied prior to averaging). Shaded regions denote the mean ±1 standard deviation.

relatively fast fall velocities (see Table 1) and the dominance of large snow hydrometeors over \( Z_{e,\text{tot,att}} \) at these levels and not the product of vertical air motion being convolved into the radar moments only in the observations.

Fig. 7 delineates the three hydrometeor classification methods currently implemented in EMC² applied over the processed model output, with the COSP emulator "observing" the model domain from the top down, similar to the on-line simulator. Congruent with the description above of the calculated radar and lidar variables, the radar-sounding and modified fixed lidar threshold methods show the domination of liquid hydrometeor classes above \( \sim 1.0 \text{ km} \), the prevalence of precipitating hydrometeors at lower levels (with lower occurrence using the radiation scheme approach), and the occasional full attenuation of the simulated lidar signal in the case of the lidar-based classification method. The COSP emulator detects a clear liquid-bearing subcolumn bin signal at cloud tops. However, because the cloud-top layer is highly reflective, generating large lidar scatter-
ing ratio values (the ratio of total to molecular attenuated backscatter), most of the underlying layers are either classified as "undefined" or generate signals too weak to be detected.

Figure 7. Hydrometeor classification of ModelE3 configuration Tun3 SCM output for 05:00 UTC processed using EMC²’s microphysics (left) and radiation (middle) approaches while considering signal attenuation and detectability: (from top to bottom) radar-sounding method, modified fixed lidar variable threshold method, and the COSP lidar simulator emulator (top-down view). (Right) Brief summary of each classification method.

Using each of the three classification methods, phase ratio statistics can be generated with EMC² offering a method for analyzing the SCM simulation. Fig. 8 portrays the temporal evolution of the SCM simulation from the view of the simulated instruments and classification methods using the microphysics or radiation approaches. When radar-sounding and fixed lidar

Radar-Sounding
(1) Cloud water > 0 → Cloud
(2) \( Z_e \geq Z_{\text{thm}} \) → Precipitation
(3) (1) + (2) → Mixed

Lidar Modified Fixed

COSP Emulator
(1) Scattering ratio (SR) > 5 → hydrometeors
(2) ATB\( \frac{\text{ATB}_{\text{ATB}}}{\text{ATB}_{\text{min}}} \) < Discrimination line\(^a\) → Liquid
(3) (2) = False → Ice
(4) Hydrometeors below (top-down) or above (bottom-up) bin with SR > 30 → Undefined

\(^a\)ATB = Attenuated total backscatter
\(^a\)Cesana and Chepfer (2013) formulation with modified coefficients implemented in ModelE3
threshold methods are applied while using the microphysics approach, the evolution of the simulated cloud appears self-consistent between the two methods and generally follows the prototypical appearance of nearly continuously precipitating liquid-bearing cloud layers with weakly varying cloud base height (e.g., de Boer et al., 2011; Fridlind and Ackerman, 2018; Silber et al., 2021a). Here, the radar-sounding skill is associated with the diminished cloud water fraction relative to the large fraction of the other hydrometeor classes (fig. 3, right) and the method’s ability to correctly detect cloud water layers. The capabilities of this method should nonetheless be considered carefully when directly compared with observations because realistic sounding profiles typically lack the fine temporal resolution emulated by EMC here, and while in-situ observations can provide a robust characterization of liquid-bearing cloud layers, they can also produce sporadic false positive or negative cloud detections (e.g., Silber et al., 2020, fig S1; Vassel et al., 2019). In the case of the modified fixed lidar threshold method, the low-level phase ratio skill originates from the hydrometeor-bearing subcolumn bins being largely classified as "ice". This classification decision is the result of the low $\beta_{p_{\text{tot}}}$ and moderate-to-high LDR (e.g., fig. 6) produced by the prevalence of ice hydrometeors relative to cloud water (fig. 3, right) and the greater mass (and likely volume due to the spherical representation) of these hydrometeors relative to rain (fig. 3, left).

The COSP emulator using the microphysics approach with a top-down view is consistent with the example in fig. 7, in which hydrometeor detection is limited by the optically-thick and highly reflective cloud-top region, resembling observational retrievals (e.g., Cesana and Chepfer, 2013; Cesana et al., 2016).

Using the radiation approach, phase ratios more frequently show sharper transitions between the extreme values (all liquid or all non-liquid) stemming from the utilization of the generalized hydrometeor fraction. In the case of the radar-sounding method, for example, there is a distinct dominance of liquid-bearing bins, which only require any amount of cloud water to be classified as such. This dominance originates in the common occurrence of some cloud water mass mixing ratios in model grid cells throughout the SCM simulation (fig. 3, left) combined with the implementation of the generalized hydrometeor fraction, which necessarily increases the overlap between cloud water and other hydrometeor classes. The limited number of model grid cells in which the subcolumn bins exhibit a more mixed behavior are the result of the randomized component of the subcolumn generator, which does not necessarily require overlap between cloud and precipitating hydrometeors (see sect. 2.2). Based on these classification results, we suggest that the radar-sounding method could lead to hydrometeor classification favoring liquid-bearing classes when the radiation approach or similar model output with a generalized hydrometeor fraction is used.

The modified fixed lidar threshold method is relatively consistent with the microphysics approach (fig. 8), even though some times (mainly around 06:00 and 10:00 UTC) are characterized by greater (smaller) relative liquid occurrence at lower (higher) model levels (phase ratio values closer to 0.5). These phase ratio differences relative to the microphysics approach are the result of the convolution of the generalized hydrometeor fraction and its deviation relative to the cloud water fraction (fig. 3, right).

Using a top-down view, the COSP emulator agrees with both its application using the microphysics approach as well as the COSP output from the on-line simulator implemented in ModelE3, which utilizes the radiation approach (fig. 8. Unlike the on-line simulator, the emulator detects some "undefined" hydrometeors at low levels (down to the surface), which can be
**Figure 8.** Time-height curtain plots showing the liquid-bearing frequency phase ratio (calculated over all hydrometeor-bearing subcolumn bins) of ModelE3 configuration Tun3 SCM simulation using the microphysics (left) and radiation (right) approaches. The classification methods used to calculate the phase ratios are (from top to bottom) the radar-sounding method, the modified fixed lidar variable threshold method, the COSP lidar simulator emulator (top-down view), and the on-line COSP lidar simulator implemented in ModelE3, which is only processed using ModelE3’s radiation approach. In this figure, the "mixed" class of the radar-sounding method is counted as liquid, while the "undefined" classes in the other two methods are treated as "non-liquid" even though in some cases they are more likely to be liquid-bearing (e.g., the "undefined1" class in the modified fixed lidar variable threshold method). The lower-left panel shows the mass phase ratio calculated using the ModelE3 output without EMC\(^2\) processing.

explained by the lack of model interpolation onto a uniform vertical grid and/or small differences in the compiled simulator code related to signal attenuation (e.g., the accumulation of optical thickness).

We note that the treatment of the COSP emulator’s "undefined" subcolumn bins as "ice" to produce phase ratio statistics leaves the impression that only ice hydrometeors exist below cloud top. However, a rather different impression of mostly liquid water dominance, though not as stark as in the radar-sounding method using the radiation approach, is perceived when the mass phase ratio calculated using the raw SCM output is examined (fig. 8, lower-left). Contrary to the COSP emulator, treating "undefined" bins as "ice" in the modified fixed lidar threshold method increases its apparent frequency phase ratio agreement with the mass phase ratio in multiple time-height bins. Phase classification depends on instrument measurement characteristics.
and limitations and hydrometeor properties such as their class, relative mixture with other hydrometeor classes, as well as their size distributions. Therefore, such an apparent agreement between different variables and phase occurrence metrics, as well as between the same variables and metrics based on different instruments and/or methodologies, should be taken with a grain of salt (cf. Cesana et al., 2021; see also Silber et al., 2021c).

Fig. 8 demonstrates the sensitivity of phase ratio statistics to the classification method, the viewing direction of the examined instrument, and the method by which "liquid" and "non-liquid" or "ice" classes are being counted. It shows that the use of forward simulators alone is not a guarantee for an "apples-to-apples" comparison, which requires matching processing steps to ensure its robustness.

4 Summary

EMC² provides an easy to use and flexible framework for the analysis of large-scale model output and its direct comparison with ground-based observations via the generation of subcolumns intended to explicitly represent a sub-grid scale variability, and the simulation of ground-based (and air- or space-borne) radars and lidars. EMC²’s framework is already tailored to the MG2 2-moment microphysics while using single-particle scattering LUTs and has the proper infrastructure required for it to be customized to other similar schemes, as well as high-resolution model output. EMC²’s option for using radiation scheme logic in the subcolumn generator and simulator enables direct comparison with other on-line active instrument simulators (e.g., the COSP lidar simulator) with a bottom-up or top-down view option that can bridge between different methodologies by evaluating differences between the outputs resulting from their implementation.

Because it is generally designed to emulate ground-based systems, EMC² is suitable for the evaluation of column output extracted from global simulations against long-term ground-based datasets. The general adaptability of the software code to other climate models and instruments via the model and instrument Python classes renders EMC² as a flexible framework to enable consistent and reproducible post-processing methods and evaluation across multiple models.

An AWARE case study was used to illustrate the application of EMC² to LES and SCM simulations of a highly super-cooled Antarctic cloud, including the utility of the program for hydrometeor classification using radar-sounding, lidar variable thresholds, and COSP emulator methods. The ModelE3 SCM using configuration Tun3 showed general agreement with the observations at the examined simulation time as well as with the baseline DHARMA LES used to develop this case study (see Silber et al., 2019a). The LES output can be processed with EMC² after a few adaptations made only to an inherited Model class (see sect. 2.1). Thus, although it was developed for large-scale models, EMC² can also be used to compare cloud resolving or LES models with observations (as shown for DHARMA). EMC² also allows the implementation of advanced scattering model calculations in the forward calculations via customized LUTs that could be matched to some scattering assumptions made by models, for example, the implementation of the MODIS C6 calculations in both ModelE3’s radiation scheme and EMC².

The AWARE case study presented here is suitable for simulation by any global model in SCM mode (see input file repository specified under code and data availability). Case study observations, as well as ModelE3 SCM and DHARMA LES inputs and
outputs from EMC\textsuperscript{2} used to produce all examples above are also provided for step-by-step illustration (see code and data availability). We plan that additional case study examples will similarly be provided to illustrate results under differing cloud regimes.

Planned future additions to EMC\textsuperscript{2} include an extension to ground-based scanning radars, a Mie scattering calculator, spectral broadening estimates for the radar simulator, and a multiple-scattering model for the lidar simulator, all of which will be configured for consistency with model physics and output fields. We invite the community to take advantage of the framework provided by EMC\textsuperscript{2} and to contribute to its further development and applications.

Appendix A: Lists of acronyms, abbreviations, and symbols
Table A1. Acronyms and subscript abbreviations used in this manuscript.

| Name     | Definition                                                                 |
|----------|---------------------------------------------------------------------------|
| ACT      | Atmospheric Community Toolkit                                             |
| ARM      | Atmospheric Radiation Measurement user facility                          |
| ATB      | attenuated total backscatter                                               |
| AWARE    | ARM West Antarctic Radiation Experiment                                    |
| CALIOP   | Cloud-Aerosol Lidar with Orthogonal Polarization                           |
| CERES    | Cloud’s and the Earth’s Radiant Energy System                             |
| COSP     | Cloud Feedback Model Intercomparison Project Observational Simulator Package|
| DHARMA   | Distributed Hydrodynamic Aerosol and Radiative Modeling Application       |
| EMC²     | Earth Model Column Collaboratory                                           |
| ESM      | Earth System Model                                                         |
| GISS     | Goddard Institute for Space Studies                                        |
| HSRL     | high spectral resolution lidar                                             |
| IWP      | ice water path (in g/m²)                                                   |
| KAZR     | Ka-band ARM zenith-pointing radar                                          |
| LDR      | linear depolarization ratio                                                |
| LES      | large eddy simulation                                                      |
| LUT      | lookup table                                                               |
| LWC      | liquid water content (in g/m³)                                             |
| LWP      | liquid water path (in g/m²)                                                |
| MG2      | Gettelman and Morrison (2015) (2-moment microphysics scheme description)  |
| MODIS    | Moderate-Resolution Imaging Spectroradiometer                              |
| MPL      | micropulse lidar                                                           |
| SCM      | single column model                                                        |
| SNR      | signal-to-noise ratio (in dB)                                              |
| WACR     | W-band ARM cloud radar                                                     |
| XSACR    | X-band scanning ARM cloud radar                                            |
| cl       | cloud water (cloud liquid)                                                 |
| ci       | cloud ice                                                                  |
| pl       | rain (precipitating liquid)                                                |
| pi       | snow (precipitating ice)                                                   |
| hyd      | a hydrometeor class (cl, ci, pl, or pi)                                     |
| tot      | a total variable, incorporating multiple hydrometeor classes and/or cloud types (convective and/or stratiform) |
| att      | attenuated (backscatter or radar reflectivity factor)                      |
| vol      | volumetric (per unit volume)                                               |
| gas      | refers to the main atmospheric gases attenuating radar signals (O₂ and H₂O) |
Table A2. Symbols of variables and parameters used in this manuscript and their units (unless explicitly stated otherwise in the text).

| Symbol     | Definition                                           | Units       |
|------------|------------------------------------------------------|-------------|
| $z$        | height                                               | meters      |
| $q$        | mixing ratio                                         | kg/kg       |
| $w$        | vertical velocity                                    | m/s         |
| $N$        | number concentration                                 | $m^{-3}$    |
| $\rho_w$  | water density (1000 kg/m$^3$)                        | kg/m$^3$    |
| $\rho_b$  | bulk density of water or ice (1000 or 917 kg/m$^3$, respectively) | kg/m$^3$ |
| $\rho_a$  | air density                                          | kg/m$^3$    |
| $\rho_{hyd}$ | density of a hydrometeor class                     | kg/m$^3$    |
| $\phi$    | particle size distribution                           | $m^{-4}$    |
| $r_e$      | effective radius                                     | meters      |
| $D$        | particle diameter                                    | meters      |
| $N_s$      | number of subcolumns                                 | -           |
| $f_{hyd}$  | fraction of a hydrometeor class                      | -           |
| $f_{gen}$  | generalized hydrometeor fraction                     | -           |
| $s, h, t$  | subcolumn (index), height (model level index), and time coordinates, respectively | -           |
| $m_{hyd}$  | complex refractive index of an hydrometeor class     | -           |
| $K_m$      | $K_m = (m_{hyd}^2 - 1)/(m_{hyd}^2 + 2)$             | -           |
| $|K_w|^2$  | dielectric factor for water                          | -           |
| $\Phi_{hyd}$ | constant fluffiness factor of an ice hydrometeor class | -           |
| $\lambda_l$ | lidar operating wavelength                          | meters      |
| $A$        | geometric cross-section                               | $m^2$       |
| $Q_e$      | extinction efficiency                                | -           |
| $Q_{bs}$   | backscattering efficiency                            | sr$^{-1}$   |
| $\beta_m$ | molecular backscatter cross-section                  | $m^{-1}sr^{-1}$ |
| $\beta_p$ | particulate backscatter cross-section                | $m^{-1}sr^{-1}$ |
| $\alpha_p$ | particulate extinction cross-section                | $m^{-1}$    |
| $T_m^2$    | two-way molecular transmittance                      | -           |
| $\tau$    | accumulated optical thickness at a model level base (bottom-up view) or top (top-down view) | -           |
| $\eta$    | multiple scattering coefficient                       | -           |
| $\lambda_r$ | radar operating wavelength                          | meters      |
| $Z_e$      | equivalent reflectivity factor                        | dBZ         |
| $V_D$      | mean Doppler velocity                                | m/s         |
| $\sigma_D$ | spectral width                                       | m/s         |
| $Z_{e,min}$ | minimum detectable $Z_e$                            | dBZ         |
| $Y$        | one-way integrated attenuation at a model level base (bottom-up view) or top (top-down view) | dB         |
Figure B1. Time series showing (top) liquid water path (LWP) and (middle) ice water path (IWP) in each case study simulation using the DHARMA LES (baseline simulation in Silber et al., 2019a), and the four ModelE3 configurations (Tun1, Tun2, Tun3, and Phys). The Eulerian retrievals shown for 10:00 UTC (error bars denote uncertainty) correspond to the time at which the Lagrangian simulated domain approaches McMurdo Station (see sect. 3.2). (Bottom, from left to right) Temperature, turbulent kinetic energy, and cloud water mixing ratio profiles at the end of the SCM and DHARMA LES simulations.

Fig B1 illustrates time series of liquid water path (LWP) and ice water path (IWP) from the ModelE3 SCM case study output using configurations Tun1, Tun2, Tun3, and Phys (see sect. 3.3), as well as the temperature, turbulent kinetic energy (TKE) and cloud water mixing ratio profiles at the end of the simulation. Out of these four model configurations, only configurations Tun3 and Phys maintain substantial amounts of LWP in this highly supercooled cloud case study and agree with both the depicted DHARMA LES output and the observed LWP retrievals (see Silber et al., 2019a). Moreover, the SCM using each of these two configurations is able to develop a cloud-top inversion and TKE, both of which are driven by radiative cooling of cloud water.
Figure B2. Mean profiles of observed (time-averaged) and simulated (subcolumn-averaged) HSRL and KAZR variables calculated using EMC\(^2\) microphysics (top) and radiation (bottom) approaches using the four ModelE3 configurations (see legend), without applying a lidar and radar extinction or SNR masks, respectively.

consistent with the DHARMA LES (see B1, bottom panels) and various polar cloud observations (e.g., Morrison et al., 2012; Silber et al., 2019a). All of ModelE3’s configurations except for Tun1, which generates the largest amount of ice, are within the range of IWP estimated retrieval uncertainty. Out of those three model configurations, Tun3 is closest to the retrieved IWP and best matches the cloud formation and evolution in the LES.

Because the supercooled cloud is developed faster in the SCM than in the LES model (see B1, top and middle), we choose to demonstrate EMC\(^2\) using the SCM output corresponding to 05:00 UTC. Fig B2 shows the mean lidar and radar variable
profiles at that model time step from EMC\textsuperscript{2} for each of the four model configurations together with the time-averaged observed profiles. Using either the microphysics or radiation approach, configurations Tun3 and Phys best match the observations. Most of the radar and lidar variables calculated using these two configurations are largely consistent with each other, but overall, the Tun3 SCM output shows the best agreement with the mean observed profiles.

To summarize, two of ModelE3’s final four configurations, that is, configurations Tun3 and Phys, show reasonable agreement with both the observed quantities as well as the LES output variables. Because configuration Tun3 performs slightly better, we focus on this model configuration for detailed comparison with LES and observations.

**Code and data availability.** The most recent EMC\textsuperscript{2} code (v. 1.2), which also supports the E3SM and CESM2 climate models as well as the Weather Research and Forecasting (WRF) model, is available on GitHub at https://github.com/columncolab/EMC2/. The EMC\textsuperscript{2} Version 1.1 code described and used in this study together with KAZR and HSRL measurements, EMC\textsuperscript{2}-processed ModelE3 and DHARMA LES model output data files, and a Jupyter Notebook demonstrating the reproduction of the plots in this manuscript using EMC\textsuperscript{2}, are available at Zenodo (http://doi.org/10.5281/zenodo.5115252; Silber et al., 2021b). The SCM initialization files (sounding + forcing) required to run the case study simulation are available at the Mendeley Data repository under http://dx.doi.org/10.17632/gz4gdn3jvjz.1.

**Author contributions.** IS and AMF conceptualized the simulator and subcolumn generator with essential support from ASA. RCJ and SC conceptualized the open-source Python framework. IS and RCJ developed the EMC\textsuperscript{2} code. ASA contributed to structuring general climate model logic and ModelE3’s logic specifics. JV supported the consistency and conceptual implementation of the instrument logic. JD provided the Yang et al. (2013) and Ding et al. (2017) single-particle scattering LUTs and added essential information required to reproduce and maintain the MODIS C6 consistency. IS prepared the manuscript with contributions from RCJ. All authors reviewed and edited the manuscript.

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References

Amante, C.: ETOPO1 1 arc-minute global relief model: procedures, data sources and analysis. NOAA Technical Memorandum NESDIS NGDC-24. National Geophysical Data Center, NOAA., Tech. rep., https://doi.org/10.7289/V5C8276M, 2009.

Bodas-Salcedo, A., Webb, M. J., Brooks, M. E., Ringer, M. A., Williams, K. D., Milton, S. F., and Wilson, D. R.: Evaluating cloud systems in the Met Office global forecast model using simulated CloudSat radar reflectivities, Journal of Geophysical Research: Atmospheres, 113, https://doi.org/10.1029/2007JD009620, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2007JD009620, 2008.

Bodas-Salcedo, A., Webb, M. J., Bony, S., Chepfer, H., Dufresne, J.-L., Klein, S. A., Zhang, Y., Marchand, R., Haynes, J. M., Pincus, R., and John, V. O.: COSP: Satellite simulation software for model assessment, Bulletin of the American Meteorological Society, 92, 1023–1043, https://doi.org/10.1175/2011BAMS2856.1, https://doi.org/10.1175/2011BAMS2856.1, 2011.

Bodas-Salcedo, A., Williams, K. D., Ringer, M. A., Beau, I., Cole, J. N. S., Dufresne, J.-L., Koshiro, T., Stevens, B., Wang, Z., and Yokohata, T.: Origins of the Solar Radiation Biases over the Southern Ocean in CFMIP2 Models, Journal of Climate, 27, 41 – 56, https://doi.org/10.1175/JCLI-D-13-00169.1, https://journals.ametsoc.org/view/journals/clim/27/1/jcli-d-13-00169.1.xml, 2014.

Bohren, C. F. and Battan, L. J.: Radar Backscattering by Inhomogeneous Precipitation Particles, Journal of the Atmospheric Sciences, 37, 1821–1827, https://doi.org/10.1175/1520-0469(1980)037<1821:RBBIPP>2.0.COhttp://0.0.0.2, 1980.

Bohren, C. F. and Huffman, D. R.: Absorption and Scattering of Light by Small Particles, John Wiley & Sons, Ltd, Weinheim, Germany, https://doi.org/10.1002/9783527618156, 1983.

Burns, D., Kollias, P., Tatarevic, A., Battaglia, A., and Tanelli, S.: The performance of the EarthCARE Cloud Profiling Radar in marine stratiform clouds, Journal of Geophysical Research: Atmospheres, 121, 14,514–525,537, https://doi.org/10.1002/2016JD025090, https://doi.org/10.1002/2016JD025090, 2016.

Campbell, J. R., Hlavka, D. L., Welton, E. J., Flynn, C. J., Turner, D. D., Spinhirne, J. D., Scott, V. S., and Hwang, I. H.: Full-Time, Eye-Safe Cloud and Aerosol Lidar Observation at Atmospheric Radiation Measurement Program Sites: Instruments and Data Processing, Journal of Atmospheric and Oceanic Technology, 19, 431–442, http://dx.doi.org/10.1175/1520-0426(2002)019{%}3C431:FSACO{%}3E2.0.CO2, 2002.

Cesana, G. and Chepfer, H.: Evaluation of the cloud thermodynamic phase in a climate model using CALIPSO-GOCCP, Journal of Geophysical Research: Atmospheres, 118, 7922–7937, https://doi.org/https://doi.org/10.1002/jgrd.50376, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/jgrd.50376, 2013.

Cesana, G. and Waliser, D. E.: Characterizing and understanding systematic biases in the vertical structure of clouds in CMIP5/CFMIP2 models, Geophysical Research Letters, 43, 10,538–10,546, https://doi.org/10.1002/2016GL070515, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2016GL070515, 2016.

Cesana, G., Kay, J. E., Chepfer, H., English, J. M., and de Boer, G.: Ubiquitous low-level liquid-containing Arctic clouds: New observations and climate model constraints from CALIPSO-GOCCP, Geophysical Research Letters, 39, https://doi.org/10.1029/2012GL053385, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2012GL053385, 2012.

Cesana, G., Chepfer, H., Winker, D., Getzewich, B., Cai, X., Jourdan, O., Mioche, G., Okamoto, H., Hagiwara, Y., Noel, V., and Reverdy, M.: Using in situ airborne measurements to evaluate three cloud phase products derived from CALIPSO, Journal of Geophysical Research: Atmospheres, 121, 5788–5808, https://doi.org/https://doi.org/10.1002/2015JD024334, https://agupubs.onlinelibrary.wiley.com/doi/ab10.1002/2015JD024334, 2016.
Cesana, G., Del Genio, A. D., Ackerman, A. S., Kelley, M., Elsaesser, G., Fridlind, A. M., Cheng, Y., and Yao, M.-S.: Evaluating models’ response of tropical low clouds to SST forcings using CALIPSO observations, Atmospheric Chemistry and Physics, 19, 2813–2832, https://doi.org/10.5194/acp-19-2813-2019, https://www.atmos-chem-phys.net/19/2813/2019/, 2019.

Cesana, G. V., Ackerman, A. S., Fridlind, A. M., Silber, I., and Kelley, M.: Snow Reconciles Observed and Simulated Phase Partitioning and Increases Cloud Feedback, Geophysical Research Letters, 48, e2021GL094876, https://doi.org/https://doi.org/10.1029/2021GL094876, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2021GL094876, e2021GL094876 2021GL094876, 2021.

Chen, Y.-S., Verlinde, J., Clothiaux, E. E., Ackerman, A. S., Fridlind, A. M., Chamecki, M., Kollias, P., Kirkpatrick, M. P., Chen, B.-C., Yu, G., and Avramov, A.: On the forward modeling of radar Doppler spectrum width from LES: Implications for model evaluation, Journal of Geophysical Research: Atmospheres, https://doi.org/10.1029/2017JD028104, http://doi.wiley.com/10.1029/2017JD028104, 2018.

Chepfer, H., Bony, S., Winker, D., Chiriaco, M., Dufresne, J.-L., and Sèze, G.: Use of CALIPSO lidar observations to evaluate the cloudiness simulated by a climate model, Geophysical Research Letters, 35, https://doi.org/https://doi.org/10.1029/2008GL034207, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2008GL034207, 2008.

Danabasoglu, G., Lamarque, J.-F., Bacmeister, J., Bailey, D. A., DuVivier, A. K., Edwards, J., Emmons, L. K., Fasullo, J., Garcia, R., Gettelman, A., Hannay, C., Holland, M. M., Large, W. G., Lauritzen, P. H., Lawrence, D. M., Lenaerts, J. T. M., Lindsay, K., Lipscomb, W. H., Mills, M. J., Neale, R., Oleson, K. W., Otto-Bliesner, B., Phillips, A. S., Sacks, W., Tilmes, S., van Kampenhout, L., Vertenstein, M., Bertini, A., Dennis, J., Deser, C., Fischer, C., Fox-Kemper, B., Kay, J. E., Kinnison, D., Kushner, P. J., Larson, V. E., Long, M. C., Mickelson, S., Moore, J. K., Nienhouse, E., Polvani, L., Rasch, P. J., and Strand, W. G.: The Community Earth System Model Version 2 (CESM2), Journal of Advances in Modeling Earth Systems, 12, e2019MS001916, https://doi.org/https://doi.org/10.1029/2019MS001916, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019MS001916, e2019MS001916 2019MS001916, 2020.

de Boer, G., Morrison, H., Shupe, M. D., and Hildner, R.: Evidence of liquid dependent ice nucleation in high-latitude stratiform clouds from surface remote sensors, Geophysical Research Letters, 38, https://doi.org/https://doi.org/10.1029/2010GL046016, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2010GL046016, 2011.

Derr, V. E., Abshire, N. L., Cupp, R. E., and McNiece, G. T.: Depolarization of Lidar Returns from Virga and Source Cloud, Journal of Applied Meteorology and Climatology, 15, 1200 – 1203, https://doi.org/10.1175/1520-0450(1976)015<1200:DOLRV>2.0.CO;2, https://journals.ametsoc.org/view/journals/apme/15/11/1520-0450_1976_015_1200_dolrv_2_0_co_2.xml, 1976.

Ding, J., Bi, L., Yang, P., Kattawar, G. W., Weng, F., Liu, Q., and Greenwald, T.: Single-scattering properties of ice particles in the microwave regime: Temperature effect on the ice refractive index with implications in remote sensing, Journal of Quantitative Spectroscopy and Radiative Transfer, 190, 26–37, https://doi.org/https://doi.org/10.1016/j.jqsrt.2016.11.026, https://www.sciencedirect.com/science/article/pii/S0022407316306367, 2017.

Doviak, R. J. and Zrnić, D. S.: Doppler Radar and Weather Observations, Academic Press, San Diego, https://doi.org/https://doi.org/10.1016/C2009-0-22358-0, 1993.

Eloranta, E. W.: Practical model for the calculation of multiply scattered lidar returns, Appl. Opt., 37, 2464–2472, https://doi.org/10.1364/AO.37.002464, http://www.osapublishing.org/ao/abstract.cfm?URI=ao-37-12-2464, 1998.

Eloranta, E. W.: High spectral resolution lidar, in: Lidar: Range-Resolved Optical Remote Sensing of the Atmosphere, pp. 143–163, Springer New York, New York, NY, 2005.

Elsaesser, G. S., Genio, A. D. D., Jiang, J. H., and van Lier-Walqui, M.: An Improved Convective Ice Parameterization for the NASA GISS Global Climate Model and Impacts on Cloud Ice Simulation, Journal of Climate, 30, 317 – 336, https://doi.org/10.1175/JCLI-D-16-0346.1, https://journals.ametsoc.org/view/journals/clim/30/1/jcli-d-16-0346.1.xml, 2017.
Falconi, M. T., von Lerber, A., Ori, D., Marzano, F. S., and Moisseev, D.: Snowfall retrieval at X, Ka and W bands: consistency of backscattering and microphysical properties using BAECC ground-based measurements, Atmospheric Measurement Techniques, 11, 3059–3079, https://doi.org/10.5194/amt-11-3059-2018, https://amt.copernicus.org/articles/11/3059/2018/, 2018.

Fan, J., Ghan, S., Ovchinnikov, M., Liu, X., Rasch, P. J., and Korolev, A.: Representation of Arctic mixed-phase clouds and the Wegener-Bergeron-Findeisen process in climate models: Perspectives from a cloud-resolving study, Journal of Geophysical Research: Atmospheres, 116, https://doi.org/https://doi.org/10.1029/2010JD015375, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2010JD015375, 2011.

Flynn, C. J., Mendozaa, A., Zhengb, Y., and Mathurb, S.: Novel polarization-sensitive micropulse lidar measurement technique, Optics Express, 15, 2785–2790, https://doi.org/10.1364/OE.15.002785, http://www.opticsexpress.org/abstract.cfm?URI=oe-15-6-2785, 2007.

Fridlind, A. M. and Ackerman, A. S.: Chapter 7 - Simulations of Arctic Mixed-Phase Boundary Layer Clouds: Advances in Understanding and Outstanding Questions, in: Mixed-Phase Clouds, edited by Andronache, C., pp. 153–183, Elsevier, https://doi.org/https://doi.org/10.1016/B978-0-12-810549-8.00007-6, https://www.sciencedirect.com/science/article/pii/B9780128105498000076, 2018.

Gettelman, A. and Morrison, H.: Advanced Two-Moment Bulk Microphysics for Global Models. Part I: Off-Line Tests and Comparison with Other Schemes, Journal of Climate, 28, 1268–1287, https://doi.org/10.1175/JCLI-D-14-00102.1, https://doi.org/10.1175/JCLI-D-14-00102.1, 2015.

Golaz, J.-C., Caldwell, P. M., Van Roekel, L. P., Petersen, M. R., Tang, Q., Wolfe, J. D., Abeshu, G., Anantharaj, V., Asay-Davis, X. S., Bader, D. C., Baldwin, S. A., Bisht, G., Bogenschutz, P. A., Branstetter, M., Brunke, M. A., Brus, S. R., Burrows, S. M., Cameron-Smith, P. J., Donahue, A. S., Deakin, M., Easter, R. C., Evans, K. J., Feng, Y., Flanner, M., Foucar, J. G., Fyke, J. G., Griffin, B. M., Hannay, C., Harrop, B. E., Hoffman, M. J., Hunke, E. C., Jacob, R. L., Jacobsen, D. W., Jeffery, N., Jones, P. W., Keen, N. D., Klein, S. A., Larson, V. E., Leung, L. R., Li, H.-Y., Lin, W., Lipscomb, W. H., Ma, P.-L., Mahajan, S., Maltrud, M. E., Mametjanov, A., McClean, J. L., McCoy, R. B., Neale, R. B., Price, S. F., Qian, Y., Rasch, P. J., Reeves Eyre, J. E. J., Riley, W. J., Ringler, T. D., Roberts, A. F., Roesler, E. L., Salinger, A. G., Shaheen, Z., Shi, X., Singh, B., Tang, J., Taylor, M. A., Thornton, P. E., Turner, A. K., Veneziani, M., Wan, H., Wang, H., Wang, S., Williams, D. N., Wolfram, P. J., Worley, P. H., Xie, S., Yang, Y., Yoon, J.-H., Zelinka, M. D., Zender, C. S., Zeng, X., Zhang, C., Zhang, K., Zhang, Y., Zheng, X., Zhou, T., and Zhu, Q.: The DOE E3SM Coupled Model Version 1: Overview and Evaluation at Standard Resolution, Journal of Advances in Modeling Earth Systems, 11, 2089–2129, https://doi.org/10.1029/2018MS001603, https://doi.org/10.1029/2018MS001603, 2019.

Hansen, J. E.: Multiple Scattering of Polarized Light in Planetary Atmospheres Part II. Sunlight Reflected by Terrestrial Water Clouds, Journal of Atmospheric Sciences, 28, 1400 – 1426, https://doi.org/10.1175/1520-0469(1971)028<1400:MSOPLI>2.0.CO;2, https://journals.ametsoc.org/view/journals/atsc/28/8/1520-0469_1971_028_1400_msopli_2_0_co_2.xml, 1971.

Heiblum, R. H., Altaratz, O., Koren, I., Feingold, G., Kostinski, A. B., Khain, A. P., Ovchinnikov, M., Fredj, E., Dagan, G., Pinto, L., Yaish, R., and Chen, Q.: Characterization of cumulus cloud fields using trajectories in the center of gravity versus water mass phase space: 2. Aerosol effects on warm convective clouds, Journal of Geophysical Research: Atmospheres, 121, 6356–6373, https://doi.org/https://doi.org/10.1002/2015JD024193, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2015JD024193, 2016.

Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas, J., Peubey, C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan, X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara, G., Dahlgren, P., Dee, D., Diamantakis, M., Dragani, R., Fleming, J., Forbes, R., Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hogan, R. J., Hólm, E., Janisková, M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay, P., Rozum, I., Vamborg, F., Vil-
laume, S., and Thépaut, J.-N.: The ERA5 global reanalysis, Quarterly Journal of the Royal Meteorological Society, 146, 1999–2049, https://doi.org/10.1002/qj.3803, https://rmets.onlinelibrary.wiley.com/doi/abs/10.1002/qj.3803, 2020.

Hillman, B. R., Marchand, R. T., and Ackerman, T. P.: Sensitivities of Simulated Satellite Views of Clouds to Subgrid-Scale Overlap and Condensate Heterogeneity, Journal of Geophysical Research: Atmospheres, 123, 7506–7529, https://doi.org/10.1029/2017JD027680, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2017JD027680, 2018.

Holz, R. E., Platnick, S., Meyer, K., Vaughan, M., Heidinger, A., Yang, P., Wind, G., Dutcher, S., Ackerman, S., Amarasinghe, N., Nagle, F., and Wang, C.: Resolving ice cloud optical thickness biases between CALIOP and MODIS using infrared retrievals, Atmospheric Chemistry and Physics, 16, 5075–5090, https://doi.org/10.5194/acp-16-5075-2016, https://acp.copernicus.org/articles/16/5075/2016/, 2016.

Hoyer, S. and Hamman, J.: xarray: N-D labeled arrays and datasets in Python, Journal of Open Research Software, 5, https://doi.org/10.5334/jors.148, http://doi.org/10.5334/jors.148, 2017.

Khairoutdinov, M. and Kogan, Y.: A New Cloud Physics Parameterization in a Large-Eddy Simulation Model of Marine Stratocumulus, Monthly Weather Review, 128, 229 – 243, https://doi.org/10.1175/1520-0493(2000)128<0229:ANCPPP>2.0.CO;2, https://journals.ametsoc.org/view/journals/mwre/128/1/1520-0493_2000_0229_ancppi_2.0.co_2.xml, 2000.

Klein, S. A. and Jakob, C.: Validation and Sensitivities of Simulated Stratocumulus Clouds Simulated by the ECMWF Model, Monthly Weather Review, 127, 2514 – 2531, https://doi.org/10.1175/1520-0493(1999)127<2514:VASOF>2.0.CO;2, https://journals.ametsoc.org/view/journals/mwre/127/10/1520-0493_1999_127_2514_vasofc_2.0.co_2.xml, 1999.

Klein, S. A., Zhang, Y., Zelinka, M. D., Pincus, R., Boyle, J., and Gleckler, P. J.: Are climate model simulations of clouds improving? An evaluation using the ISCCP simulator, Journal of Geophysical Research: Atmospheres, 118, 1329–1342, https://doi.org/10.1002/jgrd.50141, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/jgrd.50141, 2013.

Kuma, P., McDonald, A. J., Morgenstern, O., Querel, R., Silber, I., and Flynn, C. J.: Ground-based lidar processing and simulator framework for comparing models and observations (ALCF 1.0), Geoscientific Model Development, 2020, 1–45, https://doi.org/10.5194/gmd-2020-25, https://gmd.copernicus.org/articles/12/3773/2019/, 2020.

Lamer, K.: Relative Occurrence of Liquid, Ice and Mixed-phase Conditions Within Cloud and Precipitation Regimes: Long Term Ground-based Observations for GCM Model Evaluation, https://etda.libraries.psu.edu/catalog/16162kxl5431, 2019.

Lin, J.-L., Qian, T., and Shinoda, T.: Stratocumulus Clouds in Southeastern Pacific Simulated by Eight CMIP5–CFMIP Global Climate Models, Journal of Climate, 27, 3000–3022, https://doi.org/10.1175/JCLI-D-13-00376.1, https://doi.org/10.1175/JCLI-D-13-00376.1, 2014.

Lin, L., Gettelman, A., Xu, Y., Wu, C., Wang, Z., Rosenbloom, N., Bates, S. C., and Dong, W.: CAM6 simulation of mean and extreme precipitation over Asia: sensitivity to upgraded physical parameterizations and higher horizontal resolution, Geoscientific Model Development, 12, 3773–3793, https://doi.org/10.5194/gmd-12-3773-2019, https://gmd.copernicus.org/articles/12/3773/2019/, 2019.

Lubin, D., Zhang, D., Silber, I., Scott, R. C., Kalogeras, P., Battaglia, A., Bromwich, D. H., Cadeddu, M., Eloranta, E., Fridlind, A., Frossard, A., Hines, K. M., Kneifel, S., Leitch, W. R., Lin, W., Nicolas, J., Powers, H., Quinn, P. K., Rowe, P., Russell, L. M., Sharma, S., Verlinde, J., and Vogelmann, A. M.: AWARE: The Atmospheric Radiation Measurement (ARM) West Antarctic Radiation Experiment, Bulletin of the American Meteorological Society, 101, E1069–E1091, https://doi.org/10.1175/BAMS-D-18-0278.1, https://doi.org/10.1175/BAMS-D-18-0278.1, 2020.

Mätzler, C., ed.: Thermal Microwave Radiation: Applications for Remote Sensing, Electromagnetic Waves, Institution of Engineering and Technology, https://digital-library.theiet.org/content/books/ew/pbew052e, 2006.

Mech, M., Maahn, M., Kneifel, S., Ori, D., Orlandi, E., Kollias, P., Schemann, V., and Crewell, S.: PAM-TRA 1.0: the Passive and Active Microwave radiative TRAnsfer tool for simulating radiometer and radar
measurements of the cloudy atmosphere, Geoscientific Model Development, 13, 4229–4251, https://doi.org/10.5194/gmd-13-4229-2020, https://gmd.copernicus.org/articles/13/4229/2020/, 2020.

Morris, V. R.: Ceilometer Instrument Handbook, DOE/SC-ARM-TR-020, DOE Office of Science, Office of Biological and Environmental Research, 2016.

Morrison, H. and Gettelman, A.: A New Two-Moment Bulk Stratiform Cloud Microphysics Scheme in the Community Atmosphere Model, Version 3 (CAM3). Part I: Description and Numerical Tests, Journal of Climate, 21, 3642–3659, https://doi.org/10.1175/2008JCLI2105.1, https://doi.org/10.1175/2008JCLI2105.1, 2008.

Morrison, H., Thompson, G., and Tatarskii, V.: Impact of Cloud Microphysics on the Development of Trailing Stratiform Precipitation in a Simulated Squall Line: Comparison of One- and Two-Moment Schemes, Monthly Weather Review, 137, 991–1007, https://doi.org/10.1175/2008MWR2556.1, https://doi.org/10.1175/2008MWR2556.1, 2009.

Morrison, H., de Boer, G., Feingold, G., Harrington, J., Shupe, M. D., and Sulia, K.: Resilience of persistent Arctic mixed-phase clouds, Nature Geosci, 5, 11–17, http://dx.doi.org/10.1038/ngeo1332, 2012.

Myers, T. A., Scott, R. C., Zelinka, M. D., Klein, S. A., Norris, J. R., and Caldwell, P. M.: Observational constraints on low cloud feedback reduce uncertainty of climate sensitivity, Nature Climate Change, https://doi.org/10.1038/s41558-021-01039-0, https://doi.org/10.1038/s41558-021-01039-0, 2021.

Newsom, R. K.: Raman Lidar (RL) Handbook, DOE/SC-ARM-TR-038, DOE Office of Science, Office of Biological and Environmental Research, United States, https://doi.org/10.2172/1020561, https://www.osti.gov/servlets/purl/1020561, 2009.

Noel, V., Roy, G., Bissonnette, L., Chepfer, H., and Flamant, P.: Analysis of lidar measurements of ice clouds at multiple incidence angles, Geophysical Research Letters, 29, 52–1–52–4, https://doi.org/https://doi.org/10.1029/2002GL014828, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2002GL014828, 2002.

Nott, G. J. and Duck, T. J.: Lidar studies of the polar troposphere, Meteorological Applications, 18, 383–405, https://doi.org/10.1002/met.289, http://dx.doi.org/10.1002/met.289, 2011.

Oue, M., Tatarевич, A., Kollias, P., Wang, D., Yu, K., and Vogelmann, A. M.: The Cloud-resolving model Radar SIMulator (CR-SIM) Version 3.3: description and applications of a virtual observatory, Geoscientific Model Development, 13, 1975–1998, https://doi.org/10.5194/gmd-13-1975-2020, https://gmd.copernicus.org/articles/13/1975/2020/, 2020.

Penndorf, R.: Tables of the Refractive Index for Standard Air and the Rayleigh Scattering Coefficient for the Spectral Region between 0.2 and 20.0 µ and Their Application to Atmospheric Optics, Journal of the Optical Society of America, 47, 176–182, https://doi.org/10.1364/JOSA.47.000176, http://www.osapublishing.org/abstract.cfm?URI=josa-47-2-176, 1957.

Platnick, S., King, M. D., Ackerman, S. A., Menzel, W. P., Baum, B. A., Riedi, J. C., and Frey, R. A.: The MODIS cloud products: algorithms and examples from Terra, IEEE Transactions on Geoscience and Remote Sensing, 41, 459–473, https://doi.org/10.1109/TGRS.2002.808301, 2003.

Platnick, S., Meyer, K. G., King, M. D., Wind, G., Amarasinghe, N., Marchant, B., Arnold, G. T., Zhang, Z., Hubanks, P. A., Holz, R. E., Yang, P., Ridgway, W. L., and Riedi, J.: The MODIS Cloud Optical and Microphysical Products: Collection 6 Updates and Examples From Terra and Aqua, IEEE Transactions on Geoscience and Remote Sensing, 55, 502–525, https://doi.org/10.1109/TGRS.2016.2610522, 2017.

Razenkov, I. I. and Eloranta, E. W.: High spectral resolution lidar at the university of wisconsin-madison, EPJ Web Conf., 176, 01024, https://doi.org/10.1051/epjconf/201817601024, https://doi.org/10.1051/epjconf/201817601024, 2018.
Rowe, P. M., Fergoda, M., and Neshyba, S.: Temperature-Dependent Optical Properties of Liquid Water From 240 to 298 K, Journal of Geophysical Research: Atmospheres, 125, e2020JD032624, https://doi.org/10.1029/2020JD032624, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2020JD032624, e2020JD032624 2020JD032624, 2020.

Rémillard, J. and Tselioudis, G.: Cloud Regime Variability over the Azores and Its Application to Climate Model Evaluation, Journal of Climate, 28, 9707–9720, https://doi.org/10.1175/JCLI-D-15-0066.1, https://doi.org/10.1175/JCLI-D-15-0066.1, 2015.

Sassen, K.: Polarization in lidar: a review, in: Polarization Science and Remote Sensing, edited by Shaw, J. A. and Tyo, J. S., vol. 5158, pp. 151 – 160, International Society for Optics and Photonics, SPIE, https://doi.org/10.1117/12.507006, https://doi.org/10.1117/12.507006, 2003.

Segelstein, D. J.: The complex refractive index of water, Ph.D. thesis, University of Missouri–Kansas City, 1981.

Seifert, A. and Beheng, K. D.: A double-moment parameterization for simulating autoconversion, accretion and selfcollection, Atmospheric Research, 59-60, 265–281, https://doi.org/https://doi.org/10.1016/S0169-8095(01)00126-0, https://www.sciencedirect.com/science/article/pii/S0169809501001260, 13th International Conference on Clouds and Precipitation, 2001.

Shupe, M. D.: A ground-based multisensor cloud phase classifier, Geophysical Research Letters, 34, https://doi.org/https://doi.org/10.1029/2007GL031008, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2007GL031008, 2007.

Silber, I., Verlinde, J., Eloranta, E. W., and Cadeddu, M.: Antarctic cloud macrophysical, thermodynamic phase, and atmospheric inversion coupling properties at McMurdo Station. Part I: Principal data processing and climatology, Journal of Geophysical Research: Atmospheres, https://doi.org/10.1029/2018JD028279, 2018a.

Silber, I., Verlinde, J., Eloranta, E. W., Flynn, C. J., and Flynn, D. M.: Polar Liquid Cloud Base Detection Algorithms for High Spectral Resolution or Micropulse Lidar Data, Journal of Geophysical Research: Atmospheres, 123, 4310–4322, https://doi.org/https://doi.org/10.1029/2017JD027840, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2017JD027840, 2018b.

Silber, I., Fridlind, A. M., Verlinde, J., Ackerman, A. S., Chen, Y.-S., Bromwich, D. H., Wang, S.-H., Cadeddu, M., and Eloranta, E. W.: Persistent Supercooled Drizzle at Temperatures Below -25 °C Observed at McMurdo Station, Antarctica, Journal of Geophysical Research: Atmospheres, 124, 10,878–10,895, https://doi.org/10.1029/2019JD030882, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019JD030882, 2019a.

Silber, I., Verlinde, J., Wang, S.-H., Bromwich, D. H., Fridlind, A. M., Cadeddu, M., Eloranta, E. W., and Flynn, C. J.: Cloud Influence on ERA5 and AMPS Surface Downwelling Longwave Radiation Biases in West Antarctica, Journal of Climate, 32, 7935–7949, https://doi.org/10.1175/JCLI-D-19-0149.1, https://doi.org/10.1175/JCLI-D-19-0149.1, 2019b.

Silber, I., Fridlind, A. M., Verlinde, J., Russell, L. M., and Ackerman, A. S.: Nonturbulent Liquid-Bearing Polar Clouds: Observed Frequency of Occurrence and Simulated Sensitivity to Gravity Waves, Geophysical Research Letters, 47, e2020GL087099, https://doi.org/10.1029/2020GL087099, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2020GL087099, e2020GL087099 2020GL087099, 2020.

Silber, I., Fridlind, A. M., Verlinde, J., Ackerman, A. S., Cesana, G. V., and Knopf, D. A.: The prevalence of precipitation from polar supercooled clouds, Atmospheric Chemistry and Physics, 21, 3949–3971, https://doi.org/10.5194/acp-21-3949-2021, https://ACP.copernicus.org/articles/21/3949/2021/, 2021a.

Silber, I., Jackson, R. C., Fridlind, A. M., Ackerman, A. S., Collis, S., Verlinde, J., and Ding, J.: Silber_et_al_EM2_gmd-2021-194_code_and_data, https://doi.org/10.5281/zenodo.5115252, https://doi.org/10.5281/zenodo.5115252, 2021b.
Silber, I., McGlynn, P. S., Harrington, J. Y., and Verlinde, J.: Habit-Dependent Vapor Growth Modulates Arctic Supercooled Water Occurrence, Geophysical Research Letters, 48, e2021GL092767, https://doi.org/10.1029/2021GL092767, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2021GL092767, e2021GL092767 2021GL092767, 2021c.

Smith, R. N. B.: A scheme for predicting layer clouds and their water content in a general circulation model, Quarterly Journal of the Royal Meteorological Society, 116, 435–460, https://doi.org/10.1002/qj.49711649210, https://rmets.onlinelibrary.wiley.com/doi/abs/10.1002/qj.49711649210, 1990.

Sokolowsky, G. A., Clothiaux, E. E., Baggett, C. F., Lee, S., Feldstein, S. B., Eloranta, E. W., Cadeddu, M. P., Bharadwaj, N., and Johnson, K. L.: Contributions to the Surface Downwelling Longwave Irradiance during Arctic Winter at Utqiaġvik (Barrow) Alaska, Journal of Climate, https://doi.org/10.1175/JCLI-D-18-0876.1, https://doi.org/10.1175/JCLI-D-18-0876.1, 2020.

Stephens, G. L., L’Ecuyer, T., Forbes, R., Gettelman, A., Golaz, J.-C., Bodas-Salcedo, A., Suzuki, K., Gabriel, P., and Haynes, J.: Dreary state of precipitation in global models, Journal of Geophysical Research: Atmospheres, 115, https://doi.org/10.1029/2010JD014532, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2010JD014532, 2010.

Stevens, D. E., Ackerman, A. S., and Bretherton, C. S.: Effects of Domain Size and Numerical Resolution on the Simulation of Shallow Cumulus Convection, Journal of the Atmospheric Sciences, 59, 3285 – 3301, https://doi.org/10.1175/1520-0469(2002)059<3285:EODSAN>2.0.CO;2, https://journals.ametsoc.org/view/journals/atsc/59/10/jas-d-14-0265.1.xml, 2002.

Suzuki, K., Stephens, G., Bodas-Salcedo, A., Wang, M., Golaz, J.-C., Yokohata, T., and Koshiro, T.: Evaluation of the Warm Rain Formation Process in Global Models with Satellite Observations, Journal of the Atmospheric Sciences, 72, 3996 – 4014, https://doi.org/10.1175/JAS-D-14-0265.1, https://journals.ametsoc.org/view/journals/atsc/72/10/jas-d-14-0265.1.xml, 2015.

Swales, D. J., Pincus, R., and Bodas-Salcedo, A.: The Cloud Feedback Model Intercomparison Project Observational Simulator Package: Version 2, Geoscientific Model Development, 11, 77–81, https://doi.org/10.5194/gmd-11-77-2018, https://gmd.copernicus.org/articles/11/77/2018/, 2018.

Tan, I. and Storelvmo, T.: Sensitivity Study on the Influence of Cloud Microphysical Parameters on Mixed-Phase Cloud Thermodynamic Phase Partitioning in CAM5, Journal of the Atmospheric Sciences, 73, 709–728, https://doi.org/10.1175/JAS-D-15-0152.1, https://doi.org/10.1175/JAS-D-15-0152.1, 2016.

Theisen, A., Kehoe, K., Jackson, B., Sherman, Z., Godine, C., ajsockol, jkyrouac, and Hemedinger, J.: ARM-DOE/ACT: Version 0.4.3, https://doi.org/10.5281/zenodo.3987372, https://doi.org/10.5281/zenodo.3987372, 2020.

Thorsen, T. J. and Fu, Q.: Automated Retrieval of Cloud and Aerosol Properties from the ARM Raman Lidar. Part II: Extinction, Journal of Atmospheric and Oceanic Technology, 32, 1999–2023, https://doi.org/10.1175/JTECH-D-14-00178.1, https://doi.org/10.1175/JTECH-D-14-00178.1, 2015.

Tian, L. and Curry, J. A.: Cloud overlap statistics, Journal of Geophysical Research: Atmospheres, 94, 9925–9935, https://doi.org/10.1029/JD094iD07p09925, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/JD094iD07p09925, 1989.

Turner, D. D., Kneifel, S., and Cadeddu, M. P.: An Improved Liquid Water Absorption Model at Microwave Frequencies for Supercooled Liquid Water Clouds, Journal of Atmospheric and Oceanic Technology, 33, 33 – 44, https://doi.org/10.1175/JTECH-D-15-0074.1, https://journals.ametsoc.org/view/journals/atot/33/1/jtech-d-15-0074_1.xml, 2016.

Ulaby, F. T., Moore, R. K., and Fung, A. K.: Microwave Remote Sensing: Microwave remote sensing fundamentals and radiometry, Artech House, 1981.
Vassel, M., Ickes, L., Maturilli, M., and Hoose, C.: Classification of Arctic multilayer clouds using radiosonde and radar data in Svalbard, Atmospheric Chemistry and Physics, 19, 5111–5126, https://doi.org/10.5194/acp-19-5111-2019, https://acp.copernicus.org/articles/19/5111/2019/, 2019.

Wang, C., Graham, R. M., Wang, K., Gerland, S., and Granskog, M. A.: Comparison of ERA5 and ERA-Interim near-surface air temperature, snowfall and precipitation over Arctic sea ice: effects on sea ice thermodynamics and evolution, The Cryosphere, 13, 1661–1679, https://doi.org/10.5194/tc-13-1661-2019, https://tc.copernicus.org/articles/13/1661/2019/, 2019.

Warren, S. G. and Brandt, R. E.: Optical constants of ice from the ultraviolet to the microwave: A revised compilation, Journal of Geophysical Research: Atmospheres, 113, https://doi.org/https://doi.org/10.1029/2007JD009744, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2007JD009744, 2008.

Webb, M., Senior, C., Bony, S., and Morcrette, J.-J.: Combining ERBE and ISCCP data to assess clouds in the Hadley Centre, ECMWF and LMD atmospheric climate models, Climate Dynamics, 17, 905–922, https://doi.org/10.1007/s003820100157, https://doi.org/10.1007/s003820100157, 2001.

Widener, K. and Bharadwaj, N.: C-Band Scanning ARM Precipitation Radar (C-SAPR) Handbook, ARM-TR-121, DOE Office of Science, Office of Biological and Environmental Research, https://doi.org/10.2172/1054629, https://www.osti.gov/biblio/1054629, 2012.

Widener, K. and Johnson, K.: W-band ARM Cloud Radar (WACR) Handbook, ARM-TR-073, DOE Office of Science, Office of Biological and Environmental Research, United States, https://doi.org/10.2172/1019541, https://www.osti.gov/servlets/purl/1019541, 2006.

Widener, K., Bharadwaj, N., and Johnson, K.: Ka-Band ARM Zenith Radar (KAZR) Instrument Handbook, ARM-TR-106, DOE Office of Science, Office of Biological and Environmental Research, United States, https://doi.org/10.2172/1035855, http://www.osti.gov/scitech/servlets/purl/1035855, 2012a.

Widener, K. B. and Mead, J. B.: W-band ARM cloud radar—Specifications and design, in: Proc. 14th ARM Science Team Meeting, 2004.

Widener, K. B., Bharadwaj, N., and Johnson, K.: Scanning ARM Cloud Radar (X/Ka/W-SACR), ARM-TR-113, DOE Office of Science, Office of Biological and Environmental Research, United States, https://doi.org/10.2172/1043296, https://www.osti.gov/biblio/1043296-scanning-arm-cloud-radar-handbook, 2012b.

Winker, D. M.: Accounting for multiple scattering in retrievals from space lidar, in: 12th International Workshop on Lidar Multiple Scattering Experiments, edited by Werner, C., Oppel, U. G., and Rother, T., vol. 5059, pp. 128 – 139, International Society for Optics and Photonics, SPIE, https://doi.org/10.1117/12.512352, https://doi.org/10.1117/12.512352, 2003.

Winker, D. M., Vaughan, M. A., Omar, A., Hu, Y., Powell, K. A., Liu, Z., Hunt, W. H., and Young, S. A.: Overview of the CALIPSO Mission and CALIOP Data Processing Algorithms, Journal of Atmospheric and Oceanic Technology, 26, 2310 – 2323, https://doi.org/10.1175/2009JTECHA1281.1, https://journals.ametsoc.org/view/journals/atot/26/11/2009jtechal1281_1.xml, 2009.

Xiao, H., Zhao, P., Liu, X., and Li, H.: Sensitivity of aerosol-cloud interactions to autoconversion schemes in mixed-phase orographic clouds, Atmospheric Research, 247, 105 205, https://doi.org/https://doi.org/10.1016/j.atmosres.2020.105205, https://www.sciencedirect.com/science/article/pii/S0169809520311418, 2021.

Yang, P., Bi, L., Baum, B. A., Liou, K.-N., Kattawar, G. W., Mishchenko, M. I., and Cole, B.: Spectrally Consistent Scattering, Absorption, and Polarization Properties of Atmospheric Ice Crystals at Wavelengths from 0.2 to 100 um, Journal of the Atmospheric Sciences, 70, 330 – 347, https://doi.org/10.1175/JAS-D-12-039.1, https://journals.ametsoc.org/view/journals/atsc/70/1/jas-d-12-039.1.xml, 2013.

Zaremba, T. J., Rauber, R. M., McFarquhar, G. M., Hayman, M., Finlon, J. A., and Stechman, D. M.: Phase Characterization of Cold Sector Southern Ocean Cloud Tops: Results From SOCRATES, Journal of Geophysical Research: Atmospheres,
Zelinka, M. D., Myers, T. A., McCoy, D. T., Po-Chedley, S., Caldwell, P. M., Ceppi, P., Klein, S. A., and Taylor, K. E.: Causes of Higher Climate Sensitivity in CMIP6 Models, Geophysical Research Letters, 47, e2019GL085782, https://doi.org/10.1029/2019GL085782, https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019GL085782, 2020.

Zhang, Y., Xie, S., Klein, S. A., Marchand, R., Kollias, P., Clothiaux, E. E., Lin, W., Johnson, K., Swales, D., Bodas-Salcedo, A., Tang, S., Haynes, J. M., Collis, S., Jensen, M., Bharadwaj, N., Hardin, J., and Isom, B.: The ARM Cloud Radar Simulator for Global Climate Models: Bridging Field Data and Climate Models, Bulletin of the American Meteorological Society, 99, 21–26, https://doi.org/10.1175/BAMS-D-16-0258.1, 2018.