Anomalous interfacial stress generation during sodium intercalation/extraction in MoS$_2$ thin-film anodes

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Although the generation of mechanical stress in the anode material is suggested as a possible reason for electrode degradation and fading of storage capacity in batteries, only limited knowledge of the electrode stress and its evolution is available at present. Here, we show real-time monitoring of the interfacial stress of a few-layer MoS$_2$ system under the sodiation/desodiation process using microcantilever electrodes. During the first sodiation with a voltage plateau of 1.0 to 0.85 V, the MoS$_2$ exhibits a compressive stress (2.1 Nm$^{-1}$), which is substantially smaller than that measured (9.8 Nm$^{-1}$) during subsequent plateaus at 0.85 to 0.4 V due to the differential volume expansion of the MoS$_2$ film. The conversion reaction to Mo below 0.1 V generates an anomalous compressive stress of 43 Nm$^{-1}$ with detrimental effects. These results also suggest the existence of a separate discharge stage between 0.6 and 0.1 V, where the generated stress is only approximately one-third of that observed below 0.1 V. This approach can be adapted to help resolve the localized stress in a wide range of electrode materials, to gain additional insights into mechanical effects of charge storage, and for long-lifetime battery design.

INTRODUCTION

The emergence of electrical vehicles and power-hungry handheld personal data devices has stimulated an intensive research effort in high-capacity electrode materials for Li-ion batteries (LIBs) and Na-ion batteries (SIBs). Significant total capacity improvements have been achieved by using novel anode materials, such as silicon, tin, black phosphorus, and molybdenum disulfide (MoS$_2$) (1–3). Among these electrode materials, MoS$_2$ has attracted great attention due to its well-known two-dimensional (2D) structure and a large interlayer spacing of 6.15 Å (4). It has been successfully adapted as anode materials to host Li, Na, and even Mg ions (5). Its application in SIBs is of particular interest because of its high theoretical capacity of 670 mAh g$^{-1}$ (6–8).

A persistent challenge in high-capacity anode materials is the mechanical stress induced by ion intercalation/extraction, which causes large volume fluctuations and eventual failure of the electrode materials (9, 10). To understand the real-time structural changes in the electrode (11–13), various in situ technologies have been developed (14–18). For examples, in situ transmission electron microscopy (TEM) has been used to investigate the volume expansion and pulverization of single electrode particles (19–21). Similarly, atomic force microscopy (AFM) provides additional real-time mechanical information of electrodes, such as the formation of the solid-electrolyte interface layer and Young’s modulus of the electrodes (22, 23). The electrode stress state and its evolution, however, remain poorly understood. Recently, pascalammetry with operando microbattery probes was used to sense the stress in solid-state batteries (24). Although the stress developed in bulk electrodes/cell can be measured at the macroscale using techniques such as digital image correlation technology, curvature-measurement method (25, 26), and fiber optic sensors (27), a sensitive measurement modality to explicitly probe the stress at a nanoscale due to intercalation into a 2D material has yet to be proposed.

Microfabricated cantilever sensors are suspended beams that are extremely sensitive to surface stress (28). They have been successfully demonstrated as sensing elements in myriads of physical, chemical, and biological sensors that include whole-cell recognition (29), nucleic acid hybridization (30), and biomarkers and small-molecule detection (31). Microcantilevers are ideal sensors for liquid-based applications, particularly in the static mode, where dynamic coupling between the sensor and fluid is absent (32). Microcantilevers have also been used for quantitative detection of the surface stress induced by an electrochemical redox process in an electrochemistry cell (33). Here, we use this approach to investigate the stress induced during interaction/extraction of various Na ion species in few-layer MoS$_2$ under electrochemical control.

RESULTS

Multistep sodiation of MoS$_2$ in coin cells

We have used a traditional coin cell battery setup for investigating the electrochemical performance of the MoS$_2$ electrodes. We fabricated two sets of batteries using 1.7-mg ball-milled MoS$_2$ powder (see Materials and Methods). Discharging the first cell to 0.01 V shows three well-separated discharge plateaus during the first sodiation process, as shown in Fig. 1A. The first plateau occurring at 0.9 V corresponds to the intercalation of Na$^+$ to form Na$_{0.5}$MoS$_2$ according to

$$\text{MoS}_2 + x \text{Na}^+ + x e^- \leftrightarrow \text{Na}_x\text{MoS}_2 \ (x \leq 0.5) \quad (1)$$

Further intercalation of 0.5 to 1.1 mol of sodium in 1 mol of MoS$_2$ (34) following the reaction

$$\text{Na}_{0.5}\text{MoS}_2 + x \text{Na}^+ + x e^- \leftrightarrow \text{Na}_{(0.5+x)}\text{MoS}_2 \ (x \leq 0.6) \quad (2)$$

yields the second plateau at 0.8 V. The last plateau at 0.1 V is induced...
by the conversion reaction where \( \text{MoS}_2 \) is reduced to Mo and \( \text{Na}_x\text{S} \) according to

\[
\text{Na}_{1.1}\text{MoS}_2 + 2x \text{Na}^+ + 2x e^- \leftrightarrow \text{Mo} + 2\text{Na}_{(0.55+x)}\text{S}(x \leq 1.45)
\] (3)

From the second cycle onward, these plateaus blur out, and the discharge profile turns into a smooth curve. To investigate the reversibility at each stage, we subsequently cycled the second cell to four different cutoff discharge potentials (I, 1.0 V; II, 0.85 V; III, 0.4 V; and IV, 0.01 V). Above 1.0 V, the cell behaves like a highly reversible capacitor (fig. S2). The first discharge plateau at 0.9 V is reversible in all three cycles, with a corresponding charge plateau at 1.05 V (Fig. 1B). The second plateau at 0.8 V exists only in the first cycle and then turns into a smooth slope between 1.1 and 0.8 V together with the first plateau (Fig. 1C). The third discharge plateau below 0.1 V is irreversible and turns into a slope after the first cycle (Fig. 1D).

**Microcantilever platform to evaluate stress evolution**

The geometric features of a microcantilever make it an ideal substrate for the measurement of interfacial stress in the thin-film coating on the cantilever. By using a microcantilever with a thin layer of MoS\(_2\) as an electrode, it is possible to quantitatively investigate the stress evolution during the sodiation/desodiation process in the MoS\(_2\) film, as shown in Fig. 2A. To fabricate the electrode, we deposited 30 nm of silver (Ag) thin films onto both sides of the microcantilever using e-beam (electron-beam) evaporation (see Materials and Methods). As seen in the He-ion microscopy (HiM) image (Fig. 2C), a smooth silver thin film was formed on the cantilever. Later, a thin layer of MoS\(_2\) was deposited on top of the Ag-coated microcantilevers using pulsed laser deposition (PLD), as described in Materials and Methods. As shown in Fig. 2D, the surface roughness shows a slight increase with the deposited MoS\(_2\) film. The MoS\(_2\) thin film was sodiated/desodiated in a custom-made electrochemical cell with a glass window that allowed continuous monitoring of the cantilever deflection (Fig. 2B). Morphological features of the surfaces were obtained by AFM topographic measurements (Fig. 2E), which revealed that the Ag films exhibit island-like characteristics composed of a distribution of nearly spherical particles with an average size of ~20 nm. Comparatively, the MoS\(_2\) islands are significantly larger and are observed to be of quasi-hexagonal or triangular forms, which are the preferred growth structures of MoS\(_2\) (Fig. 2F). Cross-sectional high-resolution TEM was performed to assess the layered structure of the MoS\(_2\) thin film, as shown in Fig. 2G (described in Materials and Methods). The MoS\(_2\) layers are parallel to the substrate. The thickness of the thin film is ~5 nm on average, or seven to eight layers. Raman spectroscopy of the MoS\(_2\) materials, shown in Fig. 2H, demonstrates that the frequency difference between \( E_{1g} \) and \( A_{1g} \) peaks varies monotonously with the layer number of ultrathin MoS\(_2\) (35). Compared to the peaks of bulk MoS\(_2\) (the PLD target), the peaks of the MoS\(_2\) thin film shifted to higher wave numbers by ~5 cm\(^{-1}\), indicating its thickness of a few layers. Despite the marked difference in thickness, the MoS\(_2\) target and the thin film have the same crystal structure, which was confirmed by similar x-ray diffraction (XRD) patterns acquired (Fig. 2I). In both cases, the predominant (002) peak is observed, suggesting that the basal planes of the MoS\(_2\) are aligned in parallel with the testing surface and thus corroborating the thin-film cross-sectional TEM results.

To investigate the evolution of stress during the different stages, we cycled the cantilevers to the aforementioned four discharge cutoff potentials for three cycles each. Control experiments were also performed using an Ag-coated cantilever (defined as blank reference) under the same conditions to eliminate any possible environmental contributions to cantilever bending (e.g., ion adsorption). MoS\(_2\) powder (~0.1 mg) was added to the base of the cantilever array as additional electrode materials in all the tests to increase the signal/noise ratio of electrochemical signals. The correlation between the voltage and the stress is evident from Fig. 3A, displaying the simultaneous acquisition of the mechanical and electric responses. The changes in
When immersed in the electrolyte, an electrical double layer (EDL) forms at the cantilever-electrolyte interface, leading to surface stress on the blank microcantilever. In the absence of an electrochemical redox process, this stress finds its origin in both the long-range nature of the Coulombic potential and the bulk pressure of the electrolyte. The latter can be assumed to be of the same order under both conditions of electrochemical redox on MoS$_2$ and without (no change in the surface area and electrolyte dielectric properties and thus the total momentum transfer to the cantilever can be assumed the same). The Coulombic contribution to the total force (or stress component) is due to the normal component of the Maxwell stress tensor in the electrostatic limit (Supplementary Materials). The EDL-induced net stress changes nearly linearly with potential, which is consistent with the measured variations in the stress of the blank cantilever with the voltage. The net stress built up on MoS$_2$-coated cantilever is a result of the interplay between stresses induced by EDL (voltage) and those caused by sodiation/desodiation. The latter, occurring at specific potentials, can be distinguished from the EDL contribution (varying linearly with potential) by comparing with the stress changes on the blank cantilever. The results from the differential measurements are shown in Fig. 3 (B to E), where the data have been separated into four groups of curves according to the cutoff potentials. For each group, the initial stress was shifted to zero, and the total charge/discharge time was normalized to the same scale for comparison.

**Stress evolution with the cutoff potential of 1.0 V**

The stress $\Delta \sigma$ on the blank cantilever, reflecting the impact of the EDL, changes with the voltage, and the observed peaks show excellent agreement, as shown in Fig. 3B. On the MoS$_2$-coated cantilever, $\Delta \sigma$ also follows the voltage, but with noticeable differences. First, the force direction is opposite to that of the blank cantilever. As shown in Fig. 2A, both sides of the cantilever beam are coated with silver. However, because of the deep step between the base and the cantilever, the e-beam coating process deposits significantly less silver on the side wall and possibly the hinge region, leading to a higher resistance on the bottom side of the beam. This asymmetry leads to the unbalanced EDL formation on the blank cantilever, and therefore, the overall stress changes. In contrast, the MoS$_2$ coating...
(semiconductor) makes the top side less conductive than the bottom side, which changes the direction of the overall effect of EDL.

Despite the opposite directions, the stress changes of both cantilevers are expected to follow the same pattern if EDL formation is the only controlling factor. However, when the cell is discharged to near 1.0 V, a small compressive stress hump (marked as $\sigma_1$, Fig. 3B) is observed, suggesting a light sodiation process, likely involving the defects in the MoS$_2$ thin film. This process can be confirmed by the discharging curves. When the cell voltage approaches 1.0 V, the discharge rate becomes notably slower, suggesting the existence of some noncapacitive electrochemical process. In the charging process, the cantilever develops a compressive $\Delta\sigma$ up to ~2 V, which is mainly due to the formation of EDL. The overall $\Delta\sigma$ becomes tensile afterward. This compressive/tensile transition is expected to occur at the highest voltage (2.6 V) for a pure EDL-controlled process (like the blank cantilever). Apparently, there is a tensile stress from the MoS$_2$ film between 2.0 and 2.6 V (marked as $\sigma_2$), which reverses the effect of EDL. This stress may be attributed to the extraction of Na from MoS$_2$ defect sites.

**Stress evolution in the first intercalation plateau (cutoff potential of 0.85 V)**

For the group II measurements (Fig. 3C), the cell was cycled between 2.6 and 0.85 V, where the intercalation of the first 0.5 mol of Na occurs (Eq. 1). In the plateau range, the cell voltage stays almost the same, and the $\Delta\sigma$ induced by EDL remains constant, just as in the case of the blank cantilever. However, for a MoS$_2$-coated cantilever, an increasing compressive stress (marked as $\sigma_3$) was observed in the plateau range. Since the EDL contribution does not change, the observed stress reflects the stress caused by the intercalation of Na. The measured $\sigma_3$ is ~7 Nm$^{-1}$ in the first cycle and 2.1 and 1.5 Nm$^{-1}$ in the following two cycles, respectively. According to the density functional theory calculation (38) and in situ TEM analysis (39), a phase transition of MoS$_2$ from 2H to 1T polytype also occurs in this potential range in the first cycle. The stress difference of 4.9 Nm$^{-1}$ may reflect the phase transition-induced stress, which occurs only in the first cycle. The value of the stress in the second cycle serves as a better estimation of the true $\Delta\sigma$ caused by Na intercalation in the first plateau.

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**Fig. 3.** Charge/discharge curves at ~0.1 Ag$^{-1}$ (blue lines) and the synchronized stress response (black lines) of cantilever electrochemistry cells. The signals of the MoS$_2$-coated cantilever are shown in solid lines, and blank cantilever are shown in dotted lines. The cells are consequently discharged to 1.0 V (group I), 0.85 V (group II), 0.4 V (group III), and 0.01 V (group IV) for three cycles each. The information in whole test range is shown in (A), while the details in each group are shown in (B) to (E).
In the charging process (desodiation), the plateau range is not as flat as that observed during the discharging process. The EDL imparts a slowly increasing compressive stress on the MoS$_2$ cantilever (opposite to that of the blank cantilever). A tensile stress induced by Na extraction (marked as $\sigma_t$) is also expected in this range. However, this small tensile stress is buried in the compressive $\Delta \sigma$ caused by EDL. The net stress change is a small compressive stress of $1.4 \text{ Nm}^{-1}$ in the first cycle and $1.7$ and $2.2 \text{ Nm}^{-1}$ in second and third cycles, respectively. The formation of EDL is a physical process, and its impact on stress is fairly constant in different cycles. Here, the stress trend from the blank cantilevers was used as reference to estimate the stress contribution from sodiation/desodiation process. The observed increasing net compressive stress indicates that the tensile $\sigma_t$ decreases with the cycle. A similar trend was found for $\sigma_s$.

**Stress evolution in second intercalation plateau (cutoff potential of 0.4 V)**

For the group III measurements, the amount of Na intercalated into MoS$_2$ increased to 1.1 mol. In the first discharge (Fig. 3D), two separate plateaus at 0.9 and 0.8 V were observed. As can be concluded from the behavior of the blank cantilevers, the impact of EDL on $\Delta \sigma$ causes a nearly linear and very smooth slope. The stress response of MoS$_2$-coated cantilever can be separated into three sections. Besides the $\sigma_s$ in the first discharging plateau (0.9 to 0.85 V), another compressive stress (marked as $\sigma_c$) was observed in the second discharging plateau (0.8 to 0.6 V). This stress of ~3.3 $\text{N m}^{-1}$ corresponds to the Na interaction in the second plateau. Between 0.6 and 0.4 V, a larger compressive stress of 6.5 $\text{N m}^{-1}$ (marked as $\sigma_c$) was observed, indicating that the last portion of Na intercalation had a higher structural impact. From an electrochemical point of view, this result is somewhat counterintuitive. Because the discharge curve is almost linear between 0.8 and 0.4 V, the value 0.4 V has been used as a common cutoff potential in all the in situ or ex situ mechanism studies (23, 34, 38, 39). From the point of view of stress, 0.6 V is another critical point of Na intercalation.

In the charging process, the EDL initially causes a large compressive $\Delta \sigma$ with the rapid increase of the voltage. Between 1.2 and 1.5 V, $\Delta \sigma$ turns into tensile and then into compressive again. There is a tensile stress from the MoS$_2$ film (marked as $\sigma_t$) that counteracts the EDL contribution. From the charging curves (also see Fig. 1C), the range 1.2 to 1.5 V is where the main Na extraction occurs (slowest slope). $\sigma_c$ can be assigned to the tensile stress of Na extraction. In the second discharge process, the two discharge plateaus at 0.8 and 0.9 V combine into a slope at ~1 V. The corresponding $\sigma_s$ and $\sigma_c$ in the first cycle combined into a compressive stress of ~3.4 $\text{N m}^{-1}$ (marked as $\sigma_s$).

**Stress evolution in conversion reaction (cutoff potential of 0.01 V)**

When the MoS$_2$ thin film was discharged below 0.4 V (group IV measurements), a large net compressive stress of ~43 $\text{N m}^{-1}$, induced by the conversion reaction, and the EDL (tensile) effect were observed. The origin of this large stress can be explained as that from injection of up to 2.9 mol of Na into the MoS$_2$ film, which completely changes its structure (theoretically, it is Mo metal, rather than MoS$_2$). As shown in Fig. 3E, the stress change in the conversion reaction can be further separated into two sections. Before 0.1 V, the stress changes at a relatively slow rate, and the $\Delta \sigma$ is ~11 $\text{N m}^{-1}$ (marked as $\sigma_s$). The $\sigma_s$ is the extension of $\sigma_s$ observed in the group III test starting from 0.6 V. When the voltage reaches 0.1 V, the $\Delta \sigma$ slope becomes steeper, generating another additional 32 $\text{N m}^{-1}$ stress (marked as $\sigma_{10}$) before the end of the discharging process (at 0.01 V). These results suggest that the conversion reaction is composed of two separated steps with different mechanisms. A recent study based on Raman, x-ray, and TEM analyses suggests that Na$_{1.3}$MoS$_2$ formation (~0.1 V) is a critical point for the reversibility of the structure (39). Beyond that point, the MoS$_2$ is reduced into Mo metal. Our results explain the poor structural reversibility below 0.1 V as due to the generation of three times higher stress.

In the charging process, after a short initial period of compressive overall stress (EDL contribution), a large tensile stress (marked as $\sigma_{11}$) is developed, which dominates the overall process from 1.0 to 2.6 V. Because the first discharge to 0.01 V irreversibly changes the structure of MoS$_2$, the center of the sodiation-induced compressive stress shifts to a higher voltage (~0.25 V) starting from the second cycles. It also becomes smaller in the third cycle.

**DISCUSSION**

To recapitulate, in Fig. 4, we present a comparison of the main findings in our work with results from the in situ experimental and modeling studies of structural changes of MoS$_2$ reported by others. The pristine 2H-MoS$_2$ layer has a honeycomb structure with the single hexagonal unit cell of 6.39 Å in size and an interlayer spacing of 6.15 Å. Following the reaction given in Eq. 1, the MoS$_2$ accommodates ~0.5 mol of Na when sodiated to 0.85 V. A previously reported ab initio modeling study reveals that the main structural change in this process is in the axial direction, where the interlayer spacing increased to 7.42 Å, while the size of the hexagonal structure in a basal plane only changed slightly (38). In our investigation, the MoS$_2$ is grown on a silver-coated cantilever using PLD in an epitaxial manner with the basal plane parallel to the cantilever surface. The expansion in the axial direction (relative to the hexagonal structure) would not generate significant stress on the cantilever (e.g., 2.1 $\text{N m}^{-1}$ in the second cycle). In contrast, following Eq. 2, the intercalation of 0.5 to 1.1 mol of Na up to 0.4 V creates mainly a horizontal expansion in the basal plane for 2H-MoS$_2$, where the size of hexagonal structure increased to 6.58 Å, while the interlayer spacing remains almost identical (38). A phase transition from 2H to 1T MoS$_2$ is also involved in the sodiation process (39). Nevertheless, the expansion of 1T-MoS$_2$ in this potential range is also horizontal. This in-plane expansion generates a large compressive stress of 9.8 $\text{N m}^{-1}$ on the cantilever. According to in situ AFM studies, this stress is strong enough to distort the MoS$_2$ layers and induce wrinkles (23). In the conversion reaction range (0.4 to 0.01 V; Eq. 3), MoS$_2$ is converted into Mo metal particles embedded in Na$_n$S matrix, which has very limited reversibility. The conversion reaction generates a huge compressive stress of ~43 $\text{N m}^{-1}$.

Probing the stress buildup in 2D materials is extremely important for understanding the mechanisms leading to stress and thus for optimizing the battery design. For practical applications of MoS$_2$, because of its limited electrical conductivity, it is often composited with some conductive substrates. The most commonly used conductive agents are various carbons. Because of the nature of the 2D materials, the basal plane of MoS$_2$ is often parallel to the carbon substrates. In the case of our silver-coated cantilever, it can be seen from Fig. 4 that the MoS$_2$ film has adapted a similar parallel configuration.

Electrochemically induced deformation of a cantilever provides fundamental insight into the stress evolution in the MoS$_2$ thin film during the charge and discharge stages. The obtained results suggest
that by designing a proper electrochemical environment, the microcantilever measurement modality has the required sensitivity to distinguish the stress induced by the intercalation process from that of the EDL. From the point of view of stress, three discharge sections are suitable for Na storage, including the two discharge plateaus (0.9 and 0.8 V) and a slope range (0.6 to 0.1 V). Our results show that the discharge range below 0.1 V should be avoided since compared with the higher voltage range, it generates three times higher stress. Furthermore, we suggest using flexible conductive substrates with strong affinity to MoS₂ to accommodate the stress, especially in the basal plane.

A practical (bulk) electrode in batteries is composed of a mixture of active material, conductive agent, and binder. In the charge/discharge process, the stress builds up on each active material particle, while the binder and conductive agent partially accommodate the stress up to a certain extent. The macroscopic stress on bulk electrode/cell is the overall effect of the stress generation and adsorption in the electrode matrix, but it does not reflect the localized intrinsic stress on the active material particles. The localized stress is likely of more importance in understanding the electrode degradation because the localized degradation of the active material particles (e.g., cracking and pulverization) leads to battery failures. At bulk electrode/cell level, there may not be any noticeable difference between a functional and a failed cell. To develop a long-life electrode, it is essential to understand how the stress builds up at nano/microscale, e.g., in a nanofilm or in a single particle with different morphology, and under various charge/discharge conditions. The protocol developed in this work can be readily adapted to resolve the localized stress in a wide range of nano/microscale electrode materials and provide additional insights for electrode design and optimization.

**MATERIALS AND METHODS**

**Coin cell battery test**

MoS₂ powder (<2 μm) purchased from Sigma-Aldrich was first ballmilled with 20 weight % (wt %) carbon black and 10 wt % polyvinylidene...
The target used in PLD was made from the same MoS2 powder used ~1.7 mg of MoS2. The active electrodes, along with a polyethylene separator, the electrolytes, and Na counter electrodes, were assembled into a CR2032-type coin cell. The electrolyte is 1 M NaClO4 in ethylene carbonate/diethyl carbonate (1:1). All cell assemblies were operated in an argon-filled glove box. Galvanostatic charge-discharge cycling was performed on a Solartron 1470E Multichannel Potentiostat/CellTest System.

**Cantilever platform**

Microcantilevers (Octosensismicrocantilever arrays) were purchased from Micromotive GmbH. These arrays contain eight silicon (Si) microcantilevers on each chip (fig. S1). The beams are 500 µm in length, 90 µm in width, and 5 µm in thickness. Thirty nanometers of silver (Ag) was deposited on both sides of cantilevers with a 5-nm chromium (Cr) adhesion layer in a Kurt J. Lesker electron beam evaporator, serving as a current collector (top side) and a reflective layer (bottom side). The surface stress on the cantilever was calculated according to the Stoney equation, \( \Delta z = 3L^2(1 - \nu)\Delta\sigma/(E^2) \), where \( \Delta z \) is the deflection, \( L \) is the length of the cantilever beam, \( \nu \) is the Poisson’s ratio, \( \Delta\sigma \) is the surface stress, \( E \) is Young’s modulus of the cantilever, and \( r \) is the thickness of the cantilever beam. The linear stress-displacement relation presented by the Stoney equation was computationally verified for our case using the full formulation of the continuum mechanics without making any approximation. For a 3D model of the cantilever with fixed-free boundary conditions, the resulting equations were solved using the finite elements method. The result is shown in fig. S4.

**MoS2 thin-film preparation**

A thin layer of MoS2 was then deposited on the top silver layer via PLD at 500°C using ultrashort (20 ns wide) pulses from an excimer laser (KrF, \( \lambda = 248 \text{ nm} \)) with 120-mJ energy at 5-Hz repetition rate. The target used in PLD was made from the same MoS2 powder used for the coin cell battery test. PLD is based on the photon interaction between a strong laser beam (normally in the range of millijoule) and target materials to create an ejected plume, which is collected on a substrate placed at a short distance from the target (5 cm in our experiment). Unlike thermal evaporation, e-beam evaporation, or sputtering, the laser-induced expulsion produces a plume of material with stoichiometry similar to the target. It is widely used in growing epitaxial films with stoichiometry on well-matched epitaxial substrates (40). It has been reported that Ag has relative matching epitaxial parameters with MoS2 (41). The PLD target was made from the same MoS2 powder used in the coin cell battery test.

**Cantilever electrochemical cell**

It has been well understood that LIB or SIB cells are not ideal stable systems. Because the Fermi level of Li/Na is above the lowest unoccupied molecular orbital of electrolytes, the decomposition of electrolytes cannot be fully avoided (42). The small amount of impurities in electrolytes also contributes some side reactions. When the mass loading of the active electrode material is high enough (1.7 mg in the coin cell), the current from side reactions can be ignored. However, the thin MoS2 layer on one 500 µm by 90 µm cantilever is only around 1 ng; thus, its sodiation/desodiation current will be buried in the currents from side reactions. This is also a common problem for all nanobatteries used for in situ TEM and AFM studies. In this work, ~0.1 mg of additional ball-milled MoS2 powder (the same slurry as used in coin cell test) was added to the base of the cantilever array to increase the signal/noise ratio of the electrochemical current. As shown in Fig. 2A, the MoS2 powder on the array base and the thin film on cantilever were electrically connected. Both of them contribute current upon the intercalation/ extraction of Na, but only the MoS2 thin film on the cantilever beams leads to the bending of cantilevers. A small piece of Na metal connected to potassium via another platinum (Pt) wire was placed ~2 mm away from the microcantilever array and worked as the reference/counter electrode. Both electrodes were sealed in a homemade electrochemical cell (fig. 2B) filled with electrolytes. The cell had a glass window at the top, through which a focused laser beam shone on the tip of the microcantilever, and a position-sensitive detector (PSD) was used to monitor the position shifting of the reflected laser beam. The bending of the cantilever was carefully calibrated with a laser Doppler vibrometer. In our system, 1 mV of the signal on PSD corresponds to a bending of 12.3 nm on the cantilever tip, and a surface stress difference of 73.8 mNm\(^{-1}\) according to the Stoney equation. A constant current of 10 µA was used in all the charge/discharge process, which corresponds to ~0.1 A g\(^{-1}\) according to the mass of the MoS2.

**Materials characterizations**

XRD analysis of the samples was performed using a Bruker D8 Discover Diffractometer. Raman spectroscopy analysis was performed with a confocal microprobe Raman system (Thermo Nicolet Almega XR Raman microscope). Scanning HiM analysis was performed using a Zeiss Orion NanoFab equipped with Ga focused ion beam. The AFM topography was acquired using Bruker Dimension Icon system. TEM analysis was performed using a JEM-ARM200F transmission electron microscope, operating at a 200-kV accelerating voltage.

**SUPPLEMENTARY MATERIALS**

Supplementary material for this article is available at http://advances.sciencemag.org/cgi/content/full/5/1/eaav2820/DC1

Fig. S1. Scanning electron microscopy image of the eight-microcantilever array used in this work.

Fig. S2. Galvanostatic charge/discharge curves of MoS2 half-cell (versus Na) batteries at 0.05 A.

Fig. S3. TEM cross-sectional image of continuous MoS2 thin film in a long range.

Fig. S4. Linear displacement-stress relationship for a cantilever in response to a surface force in the direction perpendicular to the undisturbed surface.

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