LONG-TERM MONITORING OF THE BLACK HOLE CANDIDATE SWIFT J1753.5–0127 WITH INTEGRAL/SPI

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ABSTRACT

The black hole candidate (BHC) Swift J1753.5–0127 went into outburst in 2005 June. Rather than fade into quiescence as most BHC transients do, it has remained in a low hard spectral state for most of the ~9 years after its outburst. Its persistent emission during the hard state is reminiscent of the black hole Cyg X-1 and the BHCs 1E 1740.7–2942 and GRS 1758–258. Thus far, hard X-ray/soft gamma-ray results have mainly focused on the 2005 flare, with a few additional observations during 2007. Here, we present results from INTEGRAL/SPI observations from 2005–2010 spanning the 22–650 keV energy range. Spectral analysis shows a weak high-energy excess (∼2.9σ) above a cutoff power-law model that is well fit by a power law, suggesting an additional spectral component. Observations of Cyg X-1, 1E 1740.7–2942, and GRS 1758–258 have shown similar spectra requiring an additional high-energy component. The SPI results are compared to previously reported results for Swift J1753.5–0127 as well as observations of other sources.

Key words: black hole physics – X-rays: binaries – X-rays: individual (Swift J1753.5–0127)

1. INTRODUCTION

Swift J1753.5–0127 was discovered by the Swift/Burst Alert Telescope (BAT) on MJD 53551 (2005 June 30; Burrows et al. 2005; Palmer et al. 2005). The 15–50 keV flux continued to increase until ~ MJD 53560 (2005 July 9) with a peak flux of ~360 mCrab. Instead of returning to a quiescent state after its outburst, Swift J1753.5–0127 has maintained an average flux of ~55 mCrab in the 15–50 keV energy band. X-ray and gamma-ray observations during the flare showed spectral features and quasi-periodic oscillations consistent with a black hole (BH) in the low hard state (Miller et al. 2006; Cadolle Bel et al. 2007; Zhang et al. 2007). Consequently, Swift J1753.5–0127 is considered a black hole candidate (BHC). Radio observations during this time also detected the source, as expected for BHs in the low hard state, but the reported radio fluxes from Cadolle Bel et al. (2007) and Soleri et al. (2010) were significantly lower than is expected from radio/X-ray correlations (Gallo et al. 2003, 2006). Zuniga et al. (2008) estimated an orbital period of \( P_{\text{orb}} = 3.23 \) hr based on the period of the observed superhump while Neustroev et al. (2014) used photometric and spectroscopic variability to determine an orbital period of \( P_{\text{orb}} = 2.85 \) hr. A companion star has not yet been identified. The inability to identify it in archival images suggests that the donor star is a low-mass star (Cadolle Bel et al. 2007; Neustroev et al. 2014). Cadolle Bel et al. (2007) and Durant et al. (2009) estimate the distance to the system to be ~6–7 kpc, which suggests that the donor is a main-sequence type K or M star. Thus, the mass of the compact object is not well constrained. Froning et al. (2014) estimate a mass >10 \( M_\odot \) for a source distance of 6 kpc while Neustroev et al. (2014) report the mass as <5 \( M_\odot \).

Most BH(C)’s do not exhibit persistent or quasi-persistent emission. Remillard & McClintock (2006) list 8 such sources (BHs: Cyg X-1, LMC X-1, LMC X-3, and GRS 1915+105; BHCs: 1E 1740.7–2942, 4U 1755–388, 4U 1957+115, and GRS 1758–258) out of 20 confirmed BHs and 20 BHCs. Of the eight sources, Cyg X-1, 1E 1740.7–2942, and GRS 1758–258 are the only ones to spend most of their time in a hard state, as Swift J1753.5–0127 has done for most of the ~9 years since the flare (Soleri et al. 2013). These three sources have high-energy tails seen above ~200 keV (see Jourdain & Roques 1994 and Cadolle Bel et al. 2006 for Cyg X-1; Bouchet et al. 2009 for 1E 1740.7–2942; and Pottschmidt et al. 2008 for GRS 1758–258). Possible mechanisms for the origin of the hard tails include Comptonization of a non-thermal population of electrons (Wardzinski & Zdziarski 2001), bulk-Comptonization (Laurent & Titarchuk 1999), and jet emission (Markoff et al. 2005). INTEGRAL observations of Cyg X-1 by Laurent et al. (2011) and Jourdain et al. (2012a) have reported highly polarized emission >200 keV. The strong degree of polarization suggests the hard X-ray emission above ~200 keV is related to the radio jet seen from Cyg X-1 in the hard state.

Even though Cyg X-1, 1E 1740.7–2942, GRS 1758–258, and Swift J1753.5–0127 spend most of their time in a hard state, the observed behaviors of these sources are quite different. Historically, Cyg X-1 undergoes transitions to a soft state on the timescale of a few years, though since 2011 it has spent long periods in a soft state with short periods in the hard state. Both 1E 1740.7–2942 and GRS 1758–258 are almost always observed to be in a hard state, with short periods lasting weeks to a few months in a soft state before returning to a hard state. Swift J1753.5–0127 has exhibited extremely different behavior in that it was discovered during an X-ray outburst but did not fade back into quiescence, instead maintaining a persistent flux and remaining in a hard state for ~9 years.

In this paper, we present an analysis of INTEGRAL/SPI observations of Swift J1753.5–0127, beginning with the Target of Opportunity (ToO) during the flare in 2005 (MJD 53592) until 2010 March (MJD 55284), after which SPI observes the source for only short periods (less than ~20,000 s). The SPI results are compared to previous observations of the flare, other isolated pointed-observations of Swift J1753.5–0127 by other instruments, and other transient (GRO J0422+32) and persistent (GRS 1758–258) BH(C)s in the SPI bandpass.
regions in the ASM-SPI

5 days and an orbital inclination of 2002 October 17 into an eccentric orbit with an orbital period of 3

Swift

Bottom:

These are the first published long-term results for Swift J1753.5–0127 at hard X-ray/soft gamma-ray energies.

2. INSTRUMENT AND OBSERVATIONS

The gamma-ray observatory INTEGRAL was launched on 2002 October 17 into an eccentric orbit with an orbital period of 3 days and an orbital inclination of 51°6. These parameters were selected to minimize the amount of time spent in Earth’s radiation belt (Jensen et al. 2003). The spectrometer INTEGRAL/SPI consists of 19 hexagonal germanium detectors (GeD) configured in a hexagonal pattern with a tungsten coded mask and an active coincidence shield of 91 Bismuth Germinate (BGO) crystals to reduce the background rate (Vedrenne et al. 2003). The geometrical area of the GeD detectors is 508 cm². An individual GeD is 3.2 cm in length and 69.42 mm in height. The detectors are actively cooled to ≈85 K, resulting in an energy resolution of 2–8 keV over the 20 keV–8 MeV energy range of the instrument (Roques et al. 2003).

In order to be able to generate images of the sky, INTEGRAL follows a dithering pattern with 2° off-axis pointings from the object of interest. Each pointing lasts from 0.5–1 hr such that a complete dithering pattern is completed during an observation period. The standard dithering pattern is a rectangular pattern with 1° on-axis pointing for the source of interest and 24 off-axis pointings. There is a secondary dithering pattern which is hexagonal with 1 on-axis pointing and 6 off-axis pointings (Jensen et al. 2003).

2.1. Observations and Data Analysis

The only dedicated observation of Swift J1753.5–0127 by INTEGRAL (as of revolution 1373) was as a ToO during its flare in 2005, but because of its proximity to the Galactic Plane, the source has often been observed during the mission. Since Swift J1753.5–0127 was discovered it has been observed within 12° of the detector normal in at least one pointing (≈2000 s) in 149 INTEGRAL revolutions. For this work, only those revolutions with at least 10 pointings including Swift J1753.5–0127 within 12° of the SPI pointing direction were analyzed using the SPI Data Analysis Interface (SPIDAI)3, resulting in a total exposure time of 2.55 Ms over 53 revolutions.

The flux extraction procedure in SPIDAI performs a model fit for each energy band convolving a sky model of those sources within the field of view and a background model (based on empty field observations) with the instrument response function, which is then compared to the GeD count rates in the detector plane (see Jourdain & Roques 2009 for more detail). For spectral analysis, SPI data spanning 20–650 keV were grouped into 50 logarithmically spaced energy bins. The first two energy channels were ignored in spectral fitting because of uncertainties in the energy response (Jourdain & Roques 2009), thus reducing the energy range to 22–650 keV.

3. RESULTS

3.1. Temporal Evolution

SPI observations from 13 revolutions before the flare covering MJD 52751–53490 measured an average flux of 0.8 ± 1.2 mCrab in the 22–50 keV energy band at the location of Swift J1753.5–0127. After correcting for the positive 1 mCrab (~0.075 ct s⁻¹) bias 4 the RXTE/All-Sky Monitor (ASM) measured an average flux of 0.27 ± 0.12 mCrab in the 2–12 keV band over the period of 50088–53510 (source discovered on MJD 53520), resulting in a significance of only 2.3σ, and is thus not statistically significant. BAT did not detect the source, measuring an average flux of ~0.5 ± 0.5 mCrab in the 15–50 keV band over the period of MJD 53415–53510. Thus Swift J1753.5–0127 was not significantly detected in the ~10 years prior to the flare.

3 Publicly available interface developed at IRAP to analyze SPI data. Available at http://sigma-2.cesr.fr/integral/spidai. See the description in Burke et al. (2014).
4 See more at http://heasarc.gsfc.nasa.gov/docs/xte/asm_products.html
Figure 2. SPI 22–50 keV flux vs. 3 day ASM 2–12 keV flux corresponding to the time of the SPI observation.

Figure 1 shows the ASM 2–12 keV light curve (top panel) and the BAT 15–50 keV light curve (bottom panel) from MJD 53500–53500 (2005 May 10–2010 April 14) with 3 day averages. In the bottom panel the SPI 22–50 keV fluxes have been overplotted in arbitrary flux units with red Xs, blue squares, green triangles, and black diamonds. The different symbols denote where the revolution falls on the ASM-SPI flux plot described in Section 3.2 (Figure 2). The vertical dashed lines denote January 1 of the year listed to the right (e.g., the first dashed line corresponds to 2006 January 1).

In ~10 days, Swift J1753.5–0127 went from undetected to a peak flux of ~360 mCrab in the BAT 15–50 keV energy band and ~220 mCrab in the 2–12 keV ASM data on MJD 53560 (2005 July 9). The source flux then decays exponentially with a time constant of τ ~ 40 days in the BAT data and τ ~ 28 days in the ASM data. It is during this decay that the INTEGRAL ToO is performed, spanning MJD 53593–53595 (2005 August 11–13). The flux continues to decline until ~MJD 53700 (2005 November 26) with a flux of ~25 mCrab in the BAT data and ~10 mCrab in the ASM data. The flux slowly increases for ~1000 days with 26 SPI observations during this period.

Beginning around MJD 54636 (2008 June 19), the flux increased to >100 mCrab in the 15–50 keV band for the first time since the 2005 flare (Krimm et al. 2008) with the 2–12 keV flux increasing to ~50 mCrab. After this period, the Swift J1753.5–0127 light curve exhibits larger temporal variability in both the BAT and ASM energy ranges. The first SPI observations after the flux increase are not until approximately MJD 54900 (2009 March 10) when Swift J1753.5–0127 was observed in four revolutions over the span of ~30 days.

From about MJD 55000–55500 (2009 June–2010 October), the source enters what Soleri et al. (2013) refer to as a “failed transition,” where the spectral features soften, but the source does not make a transition to a high soft state. Swift J1753.5–0127 was in the field of view in 12 revolutions from MJD 55073–55109 (2009 August 30–October 5) for early observations during the “failed transition.”

On roughly MJD 55190 (2009 December 25), the 1.5–4 keV flux seen by the MAXI/GSC increased from ~50 to 100 mCrab over the span of several days (Negoro et al. 2009). During this period, the 4–20 keV MAXI flux remained relatively unchanged while 15–50 keV BAT flux showed a gradual decrease, suggesting a hard-to-soft state transition that lasted for ~200 days before the 15–50 keV flux recovered. The final group of SPI observations (10 revolutions) occurred after this hard-to-soft transition and covers MJD 55249–55284 (2010 February 22–March 29).

3.2. Spectral Evolution

The 22–50 keV SPI flux derived from SPIDAI and the 2–12 keV ASM average flux corresponding to the time of the revolution have been plotted in Figure 2. The fluxes separate into four distinct regions above and below 0.00065 ct cm$^{-2}$ s$^{-1}$ keV in the SPI flux and above and below 2.5 ct s$^{-1}$ in the ASM flux. Region I corresponds to the bottom left quadrant (both fluxes low), Region II corresponds to the bottom right quadrant (SPI flux low and ASM flux high), Region III corresponds to the top right quadrant (both fluxes high), and Region IV corresponds to the top left quadrant (SPI flux high and ASM flux low). In Figures 1 and 2, revolutions in Region I are marked by blue squares, Region II is marked by green triangles, Region III is marked by red Xs, and Region IV is marked by black diamonds. Region IV contains only one point, which has a negative ASM flux with a large error. This point corresponds to Revolution 848 in the SPI data, and the ASM flux is based on a single measurement with a 90 s exposure. Revolution 848 has been analyzed in Region III, as its SPI flux is consistent with that region and because revolutions 840–845, 847, and 849–852 are also in Region III. This leaves Region IV empty.

In Region I, the SPI and ASM fluxes show a roughly linear correlation while in Region II the SPI and ASM fluxes show an anticorrelation. Region III presents a different relationship between the SPI and ASM fluxes. When excluding the flare observation (the highest SPI flux), the SPI flux is roughly constant while the ASM flux varies by a factor of ~2.

The individual revolution spectra in a region were summed together to create a summed spectrum with better statistics. The total exposure is 1.69 Ms for Region I, 0.31 Ms for Region II, and 0.55 Ms for Region III, with 0.14 Ms of the Region III exposure during the flare. The summed spectrum for each region was fit in XSPEC with a power-law model and a cutoff power-law model.

Region I consists primarily of observations during the tail of the flare and the slow flux increase. For the power-law model, Region I has a best-fit spectral index of $\Gamma = 1.75 \pm 0.01$ ($\sigma$ errors) and has a $\chi^2/\nu = 2.10$ ($\nu = 46$). A cutoff power-law model better fits the data, resulting in $\chi^2/\nu = 1.01$ ($\nu = 45$), with $\Gamma = 1.45 \pm 0.05$ and $E_{\text{cut}} = 264 \pm 42$ keV (Figure 3(a)). The Region II spectrum is predominately from observations after the BAT flux decrease, but the 4–20 MAXI flux remained relatively constant. Region II is well fit by a power-law model with $\Gamma = 2.08 \pm 0.07$ and has $\chi^2/\nu = 1.15$ ($\nu = 46$) (Figure 3(b)). When fit to a cutoff power-law model, the best-fit parameters are $\Gamma = 1.95 \pm 0.21$ and $E_{\text{cut}} = 546 \pm 711$ keV with $\chi^2/\nu = 1.17$ ($\nu = 45$). The cutoff energy is not constrained and thus is not required.

Region III contains the ToO observation during the flare and observations shortly after the “failed transition” began. The SPI flux from the ToO is significantly higher than the rest of the region, and consequently that observation will be analyzed independently from the other Region III observations. When the ToO observation is fit with a power-law model, the best fit
has $\Gamma = 1.78 \pm 0.01$, with $\chi^2/\nu = 2.74$ ($\nu = 46$). The data are better fit by a cutoff power-law model with $\Gamma = 1.41 \pm 0.04$ and $E_{\text{cut}} = 213 \pm 27$ keV ($\chi^2/\nu = 0.76$ and $\nu = 45$) (Figure 3(c)). When the rest of Region III is fit with a power-law model, the best-fit photon index is $\Gamma = 1.96 \pm 0.03$, with $\chi^2/\nu = 1.69$ ($\nu = 46$). The cutoff power-law model is best fit by $\Gamma = 1.40 \pm 0.10$ and $E_{\text{cut}} = 125 \pm 21$ keV ($\chi^2/\nu = 1.06$ and $\nu = 45$) (Figure 3(d)). The spectral index is similar to Region I and the flare, but the cutoff energy has decreased.

Because of similar best-fit parameters, Regions I and III (including the flare) were combined for an average cutoff power-law spectrum. This average spectrum has best-fit parameters of $\Gamma = 1.47 \pm 0.04$ and $E_{\text{cut}} = 234 \pm 26$ keV, with $\chi^2/\nu = 1.01$ ($\nu = 45$) (Figure 4) with significant flux above 300 keV (>6$\sigma$). A weak excess above the cutoff power law ($\sim 2.9\sigma$) was present between 400 and 600 keV. To try to fit the excess, a power law was added to the cutoff power-law model. The best-fit parameters are $\Gamma_1 = 1.12 \pm 0.82$, $E_{\text{cut}} = 102 \pm 82$ keV, and $\Gamma_2 = 1.75 \pm 0.42$ with $\chi^2/\nu = 0.97$ ($\nu = 43$). The $\chi^2/\nu$ is slightly smaller compared to the cutoff power-law model, but the fit parameters are poorly constrained.

4. DISCUSSION

4.1. Comparison with Other Swift J1753.5–0127 Observations

The first SPI observations of Swift J1753.5–0127 during outburst are well fit to a cutoff power law with a low spectral index of $\Gamma = 1.41$ and a cutoff energy $\sim 200$ keV. These parameters are consistent with a BH in a hard state (Remillard & McClintock 2006). Cadolle Bel et al. (2007) fit INTEGRAL and RXTE observations from the ToO over the energy range from 3 keV to 1 MeV. The initial model used was the Comptonization model by Titarchuk (1994) convolved with an absorption model, resulting in a seed photon temperature of $kT_0 = 0.51 \pm 0.08$ keV, an electron temperature of $kT_e = 88 \pm 14$ keV, and an optical depth of $\tau = 0.67 \pm 0.14$ with large residuals <20 keV. When incorporating a reflection component to better fit the data <20 keV and performing the fit using only RXTE/Proportional Counter Array (PCA), IBIS/ISGRI, and SPI data, an acceptable $\chi^2/\nu$ was found. The best-fit parameters were a seed photon temperature of $kT_0 = 0.54_{-0.02}^{+0.02}$ keV, an electron temperature of $kT_e = 150 \pm 26$ keV, an optical depth of $\tau = 1.06 \pm 0.02$, and a reflection fraction of $\Omega/2\pi = 0.32 \pm 0.03$. The addition of the reflection component significantly increased both the electron temperature and the optical depth.

Because the residuals for the compTT fit in Cadolle Bel et al. (2007) were below the SPI energy range, an initial fit was performed without the reflection component using the seed photon temperature (0.51 keV) from Cadolle Bel et al. (2007).
The best-fit parameters for this model are an electron temperature of $kT_e = 94 \pm 24$ keV and an optical depth of $\tau = 0.74 \pm 0.22$ with $\chi^2/\nu = 0.86$ ($\nu = 45$). When the reflection component is included and the photon temperature (0.54 keV) and reflection fraction (0.32) are fixed to the values from Cadolle Bel et al. (2007), the best-fit parameters are $kT_e = 98 \pm 27$ keV and an optical depth of $\tau = 0.79 \pm 0.25$, with $\chi^2/\nu = 0.74$ ($\nu = 45$). Including the reflection component to the model improves the fit to the data compared to the compTT model alone while the parameters remain consistent. When fitting without the reflection component, the electron temperatures and optical depths are consistent with the corresponding fit in Cadolle Bel et al. (2007). When the reflection component is added to the fit, the electron temperatures are inconsistent and the optical depths are marginally consistent to the corresponding model in Cadolle Bel et al. (2007).

There have been few published spectral results at hard X-ray/soft gamma-ray energies of Swift J1753.5–0127 after the flare and none after 2007 observations by RXTE (Durant et al. 2009) and Suzaku (Reynolds et al. 2010). The RXTE (Durant et al. 2009) observations took place on MJD 54262 and 54264 (2007 June 11 and 13, respectively) using both PCA and HEXTE. The data was fit over the 2–50 keV energy range with a power law were of $\Gamma = 1.548 \pm 0.005$. The nearest SPI observations are ~80 days afterward, thus a (quasi-) simultaneous comparison cannot be made. These observations occur during Region I, so a reasonable comparison can be made with the Region I average spectrum in the overlapping energy range. In the 22–50 keV range, the SPI data are best fit to a power law with $\Gamma = 1.59 \pm 0.05$, with $\chi^2/\nu = 0.91$ ($\nu = 10$), and so are consistent with these observations.

The Suzaku observations in Reynolds et al. (2010) allow for a more direct comparison with data covering MJD 54362–54366 (2007 September 19–23). These days overlap with INTEGRAL revolutions 602 and 603, which cover MJD 54361–54366 (2007 September 18–23). Reynolds et al. (2010) report no need for a cutoff model when fitting the data over the 2–150 keV energy range. The data are best fit to a power-law model with $\Gamma = 1.619 \pm 0.003$ or $\Gamma = 1.608 \pm 0.003$ depending on the $N_H$ value used in their analysis. When the SPI data for these revolutions are fit over the 22–650 keV energy range to a power law the best-fit spectral index is significantly steeper, with $\Gamma = 1.75 \pm 0.04$ with $\chi^2/\nu = 1.19$ ($\nu = 46$). When the SPI data are fit to a cutoff power-law model, the best-fit parameters are $\Gamma = 1.35 \pm 0.13$ and $E_{\text{cut}} = 198 \pm 68$ keV, with $\chi^2/\nu = 0.97$ ($\nu = 45$). The lower $\chi^2/\nu$ indicates the presence of curvature in the spectrum even if the cutoff energy is not well constrained. However, when the energy range of the SPI data is reduced to 22–150 keV, the best-fit spectral index is $\Gamma = 1.64 \pm 0.05$ with $\chi^2/\nu = 0.76$ ($\nu = 25$), consistent with the spectral indexes from the Suzaku results.

4.2. Comparison with other BH(C)s

4.2.1. Comparison of the Flare with GRO J0422+32

The light curve of Swift J1753.5–0127 during the flare (Figure 1) displays the typical characteristics of an X-ray nova with a fast rise with a slow decline (Tanaka & Shibazaki 1996). Similar behavior was seen in the BH X-ray nova A 0620–003 (Elvis et al. 1975; Kaluzienski et al. 1977), GS 2000+251 (Tsunemi et al. 1989), GRS 1124–684 (Kitamoto et al. 1992; Ebisawa et al. 1994), and GRO J0422+32 (Harrison et al. 1994). The light curve of Swift J1753.5–0127 differs from these sources in that instead of fading into quiescence, Swift J1753.5–0127 has maintained a flux between ~50–100 mCrab in the 15–50 keV energy band for ~9 years after the initial outburst.

Swift J1753.5–0127 has some other similarities to GRO J0422+32, as both sources are located at high galactic latitude (Cadolle Bel et al. 2007; Zurita et al. 2008), have short orbital periods (<3.5 hr for Swift J1753.5–0127 and 5.1 hr for GRO J0432+22; Filippenko et al. 1995), and have been observed in the hard state during outburst. GRO J0422+32 was detected up to 600 keV with a hard tail by GRANAT/SIGMA (Roques et al. 1994), van Dijk et al. (1995) combined the SIGMA data with COMPTEL data from the Compton Gamma-ray Observatory for a spectrum spanning 35 keV–30 MeV. This joint spectrum was fit with a compTT model and gave best-fit parameters with an electron temperature of $kT_e = 100 \pm 4$ keV and an optical depth of $\tau = 1.04 \pm 0.05$ with a high-energy deviation $>300$ keV. The electron temperature for GRO J0422+32 is similar to the results from this work with compTT models with and without the reflection component (98 and 94 keV, respectively). The optical depth from van Dijk et al. (1995) is larger than the optical depths from compTT fit to Swift J1753.5–0127 (0.74 and 0.79 with and without a reflection component).

4.2.2. Comparison of the Persistent Emission with GRS 1758–258

As has been reported by Soleri et al. (2013), Swift J1753.5–0127 remained in a spectral state consistent with a BH in a low hard state for years after its initial outburst. The presence of persistent emission from BH(C)s is uncommon, as most return to quiescence within weeks to months after outburst. The flare and Regions I and III show thermal Comptonized spectra typical of BH(C)s in a low hard state and are well fit by a cutoff power-law model with a cutoff energy $>100$ keV. BH(C)s are often compared to Cyg X-1 as it is the most studied BH, and Swift J1753.5–0127 has displayed behavior similar to Cyg X-1 in its timing behavior (Soleri et al. 2013) and in its broadband spectral behavior. But the Galactic Center source GRS 1758–258 perhaps provides a closer analogue than Cyg X-1, as both Swift J1753.5–0127 and GRS 1758–258 are thought to have a low-mass companion star and to undergo mass transfer by Roche lobe overflow (Rothstein et al. 2002; Neustroev et al. 2014) while Cyg X-1 has a high-mass companion transferring mass by stellar winds (Lamers et al. 1976). Swift J1753.5–0127 and GRS 1758–258 also show radio fluxes below the expected values based on observed radio/X-ray correlations (Soleri et al. 2010).

As part of INTEGRAL’s 2003 and 2004 observation schedule, the satellite spent an extensive amount of time monitoring the Galactic Plane and the Galactic Center, allowing for numerous observations of GRS 1758–258. Pottschmidt et al. (2006) grouped INTEGRAL/ISGRI and SPI data into four “epochs,” roughly two months long each. Biweekly PCA observations were also included for an energy range of 3–500 keV. Epoch 1 took place during a “dim state” where the hard X-ray flux decreased by over an order of magnitude compared to the average hard X-ray flux. The “dim state” is characterized by the blackbody thermal spectrum of a soft state ($\Gamma = 2.29$), but the source does not transition to a high flux at soft X-rays and is similar to a “failed transition”
component is required to
the BHC GRS 1758
cutoff energies fall within the same range as expected for BH
power-law indexes than those from GRS 1758
Comptonization spectra from Regions I and III show harder
1.31 0.89
and a spectral index of
Bouchet et al.
model
describes the data than the cutoff power-law + power-law
optical depth of
is
fi
1E 1740.7
The excesses for Swift J1753.5–0127
1.50 0.25
for Epochs 2 through 9 in a hard state. The combined spectrum for Epochs
2–9 showed a statistically significant hard tail extending to
800 keV
Pottschmidt et al. (2008) find that a compTT + power-law model describes the data better than a cutoff power-law +
+ power-law model, giving best-fit parameters $kT_e = 41$ keV, an
optical depth of $\tau = 1.4$, and a spectral index of $\Gamma = 1.4$. When the combined Swift J1753.5–0127 spectrum (Figure 4)
is fit to a compTT + power-law model, the best-fit parameters are $kT_e = 39 \pm 8$ keV, an optical depth of $\tau = 1.50 \pm 0.25$,
and a spectral index of $\Gamma = 1.31 \pm 0.89$, with $\chi^2/\nu = 0.93$
($\nu = 41$). Compared to GRS 1758–258, the parameters for
Swift J1753.5–0127 are similar, even if poorly constrained, as
is the case for the spectral index. Also, this model better
describes the data than the cutoff power-law + power-law model ($\chi^2/\nu = 0.97$). As a comparison, for 1E 1740.7–2942
Bouchet et al. (2009) found a spectral index of $\Gamma = 1.9$. For
Cyg X-1, Jourdain et al. (2012b) showed that a cutoff
component is required to fit the high-energy component, which
can be described by a cutoff power-law model with $\Gamma = 1.6$
and a cutoff energy of 700 keV (Jourdain et al. 2012a).

Swift J1753.5–0127 displayed spectral behavior similar to
the BHC GRS 1758–258 in both a hard spectral state and in a
“failed transition” state, and also displayed a high-energy
excess, though it was weak in the case of Swift J1753.5–0127.
The excesses for Swift J1753.5–0127 and GRS 1758–258 are
well described by hard power-law indexes $\Gamma = 1.3$–1.4 while
1E 1740.7–2942 is well fit by a steeper index, and Cyg X-1
requires a cutoff model for the high-energy component.

5. CONCLUSION

In this work, we presented the first published hard X-ray/soft
gamma-ray observations of Swift J1753.5–0127 since 2007
observations with RXTE (Durant et al. 2009) and Suzaku
(Reynolds et al. 2010). The INTEGRAL/SPI observations span
the decreasing portion of the flare in 2005 (MJD 53592) until
2010 March (MJD 55284). Because of SPI’s sensitivity at high
energies, we have been able to detect the high-energy cutoff of
Swift J1753.5–0127 that had not been seen since the flare. The
source remained in a persistent, low hard flux state for a
majority of the SPI observations, undergoing a “failed
transition” (Soleri et al. 2013) beginning in 2009 June
(~MJD 54983).

An analysis of SPI data for 53 INTEGRAL revolutions, when
combined with RXTE/ASM data, revealed three distinct flux
regions after the initial outburst. For each region, the summed
spectrum was fit to a power-law model and a cutoff power-law
model over the 22–650 keV energy range. Region I (when both
SPI and ASM fluxes were low) was well fit by a cutoff power-law
model with a spectral index of 1.45 and a cutoff energy of
265 keV. Region II (when the SPI flux was low and the ASM
flux was high) was well fit by a power-law model with a
spectral index of 2.08. Region III (when the SPI flux was
roughly constant and the ASM flux varied by ~2) was well fit
by a cutoff power-law model with a spectral index of 1.40 and
1.54 1.69
0127 shows an excess as having a similar
spectral behavior similar to
GRS 258
265 keV. Region II
spectral similarity between the electron temperatures and
optical depths of both sources. The continued emission for
four sources have
39 keV, an optical depth of $\tau = 1.50$, and a spectral index of $\Gamma = 1.31$.

In conclusion, Swift J1753.5–0127 presents the unusual case of
a BHC initially behaving like a typical X-ray nova until failing
to fade into quiescence. The source showed some
temporal similarities to GRO J0422+32 during the flare,
with a fast rise time and a slow decay, as well as some
spectral similarities between the electron temperatures and
optical depths of both sources. The continued emission for
Swift J1753.5–0127 while in a hard spectral state puts the
source in a small group of BH(C)s consisting of Cyg X-1, 1E 1740.7–2942, and GRS 1758–258. These four sources have
(quasi-)persistent emission and remain in a hard state most
of the time. Above ~200 keV, the spectra of Cyg X-1, 1E 1740.7–2942, and GRS 1758–258 show an additional spectral component. Swift J1753.5–0127 shows a weak excess above ~200 keV but is not as bright as Cyg X-1 and lacks the exposure time of 1E 1740.7–2942 and GRS 1758–258 (~8 Ms of
INTEGRAL observations).

For Cyg X-1, polarization measurements suggest that the
hard tail is associated with a jet. As Cyg X-1, 1E 1740.7–2942,
and GRS 1758–258 are all microquasars (Mirbel et al. 1992; Rodriguez et al. 1992; Stirling et al. 1992), the hard tails for
1E 1740.7–2942 and GRS 1758–258 could also be related to
jets. Interestingly, Soleri et al. (2010) report 1E 1740.7–2942,
GRS 1758–258, and Swift J1753.5–0127 as having a similar
radio/X-ray relationship (with Cyg X-1 not reported). Assuming
the persistent emission continues, Swift J1753.5–0.127, as
well as 1E 1740.7–2942 and GRS 1758–258, are candidates for
polarization above the thermal Comptonization component
with a future mission that has higher polarization sensitivity.

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