A five-dimensional Riemannian manifold with an irreducible SO(3)-structure as a model of abstract statistical manifold

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Abstract In the present paper, we consider a five-dimensional Riemannian manifold with an irreducible SO(3)-structure as an example of an abstract statistical manifold. We prove that if a five-dimensional Riemannian manifold with an irreducible SO(3)-structure is a statistical manifold of constant curvature, then the metric of the Riemannian manifold is an Einstein metric. In addition, we show that a five-dimensional Euclidean sphere with an irreducible SO(3)-structure cannot be a conjugate symmetric statistical manifold. Finally, we show some results for a five-dimensional Riemannian manifold with a nearly integrable SO(3)-structure. For example, we prove that the structure tensor of a nearly integrable SO(3)-structure on a five-dimensional Riemannian manifold is a harmonic symmetric tensor and it defines the first integral of third order of the equations of geodesics. Moreover, we consider some topological properties of five-dimensional compact and conformally flat Riemannian manifolds with irreducible SO(3)-structure.

Keywords Information geometry · Statistical manifold · Riemannian manifold · SO(3)-structure

Mathematics Subject Classification (2010) 53C20 · 53C21 · 53C24

1 Introduction

In the present paper, we consider a five-dimensional Riemannian manifold with an irreducible SO(3)-structure as an example of an abstract statistical manifold used in information geometry. On one side, our paper shows new geometrical interpretation to the theory of abstract...
On the other side, we study new geometrical and topological properties of five-dimensional manifolds with SO(3)-structures. This paper is a continuation of studies of abstract statistical manifolds, which the second author began in her papers [49,50].

The paper is organized as follows. In Sect. 2, we give brief overview of the information geometry, the theories of abstract statistical manifolds and irreducible SO(3)-structures. In Sect. 3 of the paper we consider conjugate connections on statistical manifolds which are one of the fundamental concepts of information geometry. We also consider properties of a well-known conjugate symmetric statistical manifold and define a harmonic statistical manifold. In addition, we prove “vanishing theorems” about obstructions to the existence of compact conjugate symmetric statistical manifolds with the positive definite curvature operator and compact harmonic statistical manifolds with the negative definite curvature operator. For this, we use the Bochner technique that is one of the oldest and most important techniques in modern Riemannian geometry [41, pp. 187–234]. In Sect. 4, we study a five-dimensional Riemannian manifold with an irreducible SO(3)-structure as a model of an abstract statistical manifold. For example, we prove that if a five-dimensional Riemannian manifold with an irreducible SO(3)-structure is a statistical manifold of constant curvature, then the metric of this Riemannian manifold is an Einstein metric. In addition, we show that a five-dimensional Euclidean sphere with an irreducible SO(3)-structure cannot be a conjugate symmetric statistical manifold. Moreover, we consider some topological properties of such manifolds. Finally, we present some results for a five-dimensional Riemannian manifold with a nearly integrable SO(3)-structure. For example, we prove that the structure tensor of a nearly integrable SO(3)-structure on a five-dimensional Riemannian manifold is a harmonic symmetric tensor and it defines the first integral of third order of the equations of geodesics. Moreover, we show that an irreducible SO(3)-structure on a five-dimensional Riemannian manifold cannot be nearly integrable if the curvature operator of Riemannian manifold is negative definite.

2 Basis definitions and notations

2.1 Information geometry

In this section, we present brief historical overview of information geometry. We want to highlight how many researchers have contributed to the development of information geometry, its theory and applications.

Information geometry studies invariant properties of a family of probability distributions and can be applied to various problems in science. Statisticians use statistical models to derive inferences; they use families of probability distributions which form, in most cases, a finite-dimensional manifold which in information geometry is called a statistical manifold. Then one can ask: What are the intrinsic properties of such manifold and how the geometrical structure of such manifold is related to characteristics of statistical inference? Information geometry emerged from these questions [4].

Many authors contributed to the development of information geometry and geometrical theory of statistics. The results of their research are reflected in numerous papers and the following monographs ([1,3,7,8,12,27,33,34,53] and etc.). Moreover, a number of international workshops and symposiums on this subject were held: in UK (London, July 10–14, 2000), Italy (Pescara, Sept. 1–5, 2002), Denmark (Leipzig, August 27, 2003), Canada (Toronto, May 8–18, 2004), USA (Ann Arbor, July 29, 2004), Japan (Tokyo, December
There are two main natural geometrical objects on any statistical model equipped with a differentiable structure. Namely, they are the Fisher information tensor and the Chentsov–Amari connection.

The information tensor was introduced by R.A. Fisher in 1925 as an information characterization of a statistical model [22]. In the 1945 paper [43], Rao pointed out that the Fisher information tensor determines a Riemannian metric (which is now called the information metric or Fisher metric) for the manifold obtained from the family of probability density functions.

Let us consider the recent concept of this metric in more details. Let $\Theta \subseteq \mathbb{R}^n$ be the parameter space of an $n$-dimensional smooth family defined on some fixed event space $\Omega$, i.e., $\{p_\theta | \theta \in \Theta\}$ with $\int_\Omega p_\theta = 1$ for all $\theta \in \Theta$. Under certain assumptions, the Fisher information matrix defines the unique Riemannian metric $g$ on $\Omega$ with the following components [3, p. 32], [7, p. 28]:

$$g_{ij} = E_\theta [\theta_i \log p_\theta \theta_j \log p_\theta] = \int_\Omega p_\theta \left(\frac{\partial l}{\partial \theta^i} \frac{\partial l}{\partial \theta^j}\right) = -\int_\Omega p_\theta \left(\frac{\partial^2 l}{\partial \theta^i \partial \theta^j}\right),$$

where $E_\theta$ denotes the expectation with respect to $p_\theta$, i.e., $E_\theta[f] := \int_\Omega f(x) p_\theta$, $l = \log p_\theta$ and $\theta^i$ are coordinates of $\theta \in \Theta \subseteq \mathbb{R}^n$. Then the set $S = \{p_\theta | \theta \in \Theta \subseteq \mathbb{R}^n\}$ with the Fisher metric $g$ can be considered as a Riemannian manifold $(S, g)$ [3, p. 34], [7, p. 32].

Chentsov proved that the Fisher metric is a unique invariant metric [14]. Moreover, he defined in [14] a one-parameter group of invariant affine connections in the space of statistic distributions (which is now called the Chentsov–Amari connections), leaving Efron [19] to expose the relationship between statistical curvature and the characteristics of inference. S.-I. Amari and H. Nagaoka introduced a conjugate structure (or duality structures) in information geometry, a finding that has played a fundamental role in development of more applications of information geometry [7]. In particular, the notion of dually flat metrics was first introduced by Amari and Nagaoka [7] when they studied information geometry on Riemannian spaces. Later on, Shen extended the notion of locally dually flatness for Finsler metrics [47]. He identified and studied the dually flat Finsler metrics that are a special and valuable class of Finsler metrics in Finsler geometry, which play a very important role in studying flat Finsler information structures [16].

Next, we shall inform the reader about the basic facts of the geometry of Chentsov–Amari connections. Namely, consider for $\alpha \in \mathbb{R}$ functions $^{\alpha}_{ij,k}$ which map each point $\theta \in \Theta$ to the following value [3, p. 34], [7, p. 32]:

$$^{\alpha}_{ij,k} = E_\theta \left[\left(\partial_i \partial_j l + \frac{1 - \alpha}{2} \partial_i l \partial_j l\right) (\partial_k l)\right] = \int_\Omega \left(\partial_i \partial_j l + \frac{1 - \alpha}{2} \partial_i l \partial_j l\right) \partial_k l \ p_\theta.$$

For every $\alpha \in \mathbb{R}$, the $n^3$ functions $^{\alpha}_{ij,k}$ are the Christoffel symbols of an affine $\alpha$-connection $^\alpha$ on $S$. This connection is clearly a symmetric connection. Moreover, the relationship between the $\alpha$-connection $^\alpha$ and $\beta$-connection $^\beta$ is given by (see [7, p. 33])

$$^\beta_{ij,k} - ^\alpha_{ij,k} = \frac{\alpha - \beta}{2} T_{ijk},$$
where \( T_{ijk} \) are components of a covariant completely symmetric tensor field of order 3 which is defined by the equation \( T_{ijk} = E_\theta \partial_i l \partial_j l \partial_k l \). In addition, we also note that \( \bar{\nabla} = \frac{1}{1+\alpha} \nabla + \frac{1-\alpha}{2} \nabla \) and the 0-connection is the Levi-Civita connection with respect to the Fisher metric [7, p. 33]. Numerous articles develop further the theory of the Chentsov–Amari connection.

In addition, now we can see a new trend of the information geometry which is called Quantum information theory [12,21,23,37,42]. This theory is based on information geometry and quantum mechanics. For example, Nagaoka [37] and Petz [42] studied the information geometry of quantum probability. On the other hand, Gibilisco and Isola in [21] gave a mathematical foundation that extended information geometry to the function space. In particular, they showed that a family of inequalities, which relates to the uncertainty principle, has a geometric interpretation in terms of quantum Fisher information. At the same time, Cafaro [12] considered problems of classical and quantum chaos using the information geometry.

Further on the role of differential geometry and its applications in the probability theory and statistics, one can refer to [11] and to the special issue of AISM 59 : 1 (2007) with the title “Information Geometry and its Applications”. In particular, some authors [6,51,54] reported applications of information geometry to Bayesian statistics [32] and Bartlett correction [15]. In addition, we must say that information geometry is applied not just to statistics and information theory but also to combinatorics of neural networks and psychology [5], physics [20] and many other fields.

2.2 An abstract statistical manifold

Lauritzen [8, pp. 165–215] gave a modern differential geometric treatment of statistical problems by introducing the notion of an abstract statistical manifold. This generalization has been considered in many papers by other geometers ([6,24,29,30,49,50] and etc.). In this section, we give an overview of the geometry of such manifold.

According to Lauritzen [8, p. 179], a statistical manifold is a triplet \((M, g, T)\), where \( M \) is a connected \( C^\infty \)-manifold, \( g \) is a Riemannian metric with components \( g_{ij} \) and \( T \) is a smooth covariant completely symmetric tensor field of order 3 with components \( T_{ijk} = T(\partial_i \partial_j \partial_k) \) on each chart \( U(x^1, \ldots, x^n) \subset M \), i.e., \( T \in C^\infty S^3 M \). This tensor \( T \) is called the skewness tensor [8, p. 179]. Further, Lauritzen defines a one-parameter group of affine torsion-free \( \alpha \)-connections \( \alpha \nabla \) whose Christoffel symbols \( \Gamma_{ij}^k \) on each chart \( U(x^1, \ldots, x^n) \subset M \) are related to the Cristoffel symbols \( \Gamma_{ij}^k \) of the Levi-Civita connection \( \nabla \) by the following identities:

\[
\Gamma_{ij}^k = \Gamma_{ij}^k - \frac{\alpha}{2} T_{ij}^k
\] (2.1)

for a real parameter \( \alpha \) and

\[
T_{ij}^l = g^{lk} T_{ijk},
\] (2.2)

where the inverse of the matrix \( (g_{ij}) \) is denoted by \( (g^{ij}) \). After that, Lauritzen proves the identity [8, p. 180]

\[
\alpha \nabla_k g_{ij} = \alpha \cdot T_{kij},
\] (2.3)

and hence \( \alpha \nabla \) is a covariant completely symmetric tensor field of order 3. The identities (2.3) are called the Codazzi equations and \( g \) is called a Codazzi tensor with respect to \( \alpha \)-connections \( \alpha \nabla \) [39, p. 21], [48, pp. 56, 68, 142].
It is well known from [17], [48, p. 53] that the affine torsion-free connections $\Gamma^\alpha_{ij}$ and $\Gamma^\beta_{ij}$ are called \textit{conjugate connections} (or \textit{dual connections}) relative to the non-degenerate bilinear form $g$ if their Christoffel symbols $\Gamma^\alpha_{ik}$ and $\Gamma^\beta_{ik}$ on each chart $U(x^1, \ldots, x^n) \subset M$ satisfy the following conditions:

$$\partial_k g_{ij} = g_{jl} \partial_i \Gamma^\alpha_{ik} + g_{il} \partial_j \Gamma^\beta_{ik}. \quad (2.4)$$

In this case, we say that $\{\nabla, g, \tilde{\nabla}\}$ are conjugates. In particular, from (2.4) we conclude that $\{\nabla, g, \nabla\}$ are conjugate and the Levi-Civita connection $\nabla := \nabla$ is a self-conjugate connection [8, p. 181]. Conjugate connections have a simple interpretation [39, p. 21].

Next, Lauritzen [8, p. 185] introduces a tensor field $F \in C^\infty S^3 M$ which plays an important role for the statistical manifold. It is defined as follows:

$$F_{ijkl} = \nabla_l T_{ij}. \quad (2.5)$$

If $F \in C^\infty S^4 M$, i.e., $F$ is completely symmetric, then $(M, g, T)$ is said to be a \textit{conjugate symmetric} statistical manifold. For this manifold, a simple calculation shows that

$$0 = \nabla_i \nabla_k g_{ij} - \nabla_i \nabla_k g_{ij} = -\alpha R_{ijk} - \alpha R_{ijk},$$

where $\alpha R_{ijk} = g_{jm} \alpha R_{ilk}^m$ for the local components $\alpha R_{ilk}^m$ of the $\alpha$-curvature tensor $\alpha R$ of $\nabla$. Therefore, a \textit{conjugate symmetric} statistical manifold has the following curvature: properties

$$\alpha R(X, Y, Z, V) = -\alpha R(X, Y, Z, V); \quad \alpha R(X, Y, Z, V) = \alpha R(Z, V, X, Y), \quad (2.6)$$

where $\alpha R(X, Y, Z, V) = g(\alpha R(X, Y)Z, V)$ for any vectors $X, Y, Z, V \in T_x M$ at each point $x \in M$.

On the other hand, Takeuchi and Amari showed in [52] that the sufficient condition for a statistical manifold to be conjugate symmetric is that its connection $\alpha \nabla$ is \textit{equiaffine}. If we recall from [39, p. 14], [48, pp. 57–58] that a torsion-free affine connection $\nabla$ is called an \textit{equiaffine} or \textit{Ricci-symmetric} connection if $\alpha \nabla \omega = 0$ for the volume form $\alpha \omega : x \in M \rightarrow \omega_x = \omega(X_1, X_2, \ldots, X_n)$, where $X_1, X_2, \ldots, X_n \in T_x M$ at each point $x \in M$, then we see that the Ricci tensor $\alpha \text{Ric}$ of $\nabla$ is symmetric. The converse is true.

Next, for $\alpha = 1$ the $\alpha$-curvature tensor $\alpha \text{Ric}$ and Ricci $\alpha$-curvature tensor $\tilde{\alpha} \text{Ric}$ of the connection $\tilde{\nabla}$ are called the \textit{curvature} and \textit{Ricci tensors of a statistical manifold} $(M, g, T)$, respectively [24,29,30]. The scalar $\alpha$-curvature $\tilde{s} := \text{trace}_g \alpha \text{Ric}$ is called a \textit{scalar curvature} of a statistical manifold $(M, g, T)$ if $\alpha = 1$. Moreover, the statistical manifold $(M, g, T)$ is said to be of \textit{constant curvature} if

$$\frac{1}{n(n - 1)} \tilde{s} \cdot (g_{ij} \delta^i_k - g_{jk} \delta^i_j), \quad (2.7)$$

where $\frac{1}{n} R(\partial_k, \partial_l) \partial_j = \frac{1}{n} R_{ijkl} \partial_i$ for $\partial_k = \partial / \partial x^k$. In particular, a statistical manifold $(M, g, T)$ is said to be \textit{locally flat} if $\frac{1}{n} R = 0$. 

\begin{figure}
\includegraphics{figure.png}
\end{figure}
Finally we introduce the new class of statistical manifolds that we shall call as harmonic statistical manifolds. Before we present two differential operators of order 1. Namely, the operator $\delta^* : C^\infty S^3 M \to C^\infty S^4 M$ such that [9, p. 35]; [45]

$$(\delta^* T)_{kij} = \nabla_k T_{ijl} + \nabla_i T_{jlk} + \nabla_j T_{lki} + \nabla_l T_{kij}$$

and its formal adjoint differential operator $\delta : C^\infty S^4 M \to C^\infty S^3 M$, which is defined by

$$(\delta T)_{jk} = -\nabla_l T^l_{jk}.$$ 

Then we can present the Laplacian $\Delta_{\text{sym}} : C^\infty S^3 M \to C^\infty S^3 M$ acting on smooth covariant symmetric tensor fields which is defined as $\Delta_{\text{sym}} = \delta^* \delta - \delta \delta^*$ [9, pp. 52–54], [45]. The operator $\Delta_{\text{sym}}$ is related to a variational problem as follows: If we define the “energy” of symmetric tensor field $T$ by $E(T) = \frac{1}{2} \langle \Delta_{\text{sym}} T, T \rangle$, then $\Delta_{\text{sym}} T = 0$ is the condition for a free extremal of $E(T) = \frac{1}{2} \langle T, \Delta_{\text{sym}} T \rangle$. Such tensor is called harmonic [45]. In accordance with this definition, we shall call $(M, g, T)$ a harmonic statistical manifold if $\Delta_{\text{sym}} T = 0$.

In the Sect. 4.2, we shall consider a five-dimensional Riemannian manifold $(M, g)$ with a nearly integrable SO(3)-structure as an example of a harmonic statistical manifold.

2.3 Overview of the theory of irreducible SO(3)-structures

It is well known that a Riemannian metric $g$ on an $n$-dimensional orientable manifold $M$ determines $SO(n)$-structures, where the tangent space $T_x M$ at each point $x \in M$ behaves as a representation for $SO(n)$. An arbitrary subgroup $G \subset SO(n)$ determines restricted Riemannian $G$-structure, i.e., the tangent space $T_x M$ must behave as a representation for $G$. Usually, the tangent space $T_x M$ is regarded as an irreducible representation of a subgroup $G \subset SO(n)$. Following series of papers, some [2,10,13] consider an irreducible SO(3)-structure on a five-dimensional Riemannian manifold $(M, g)$. In particular, it was shown that an irreducible SO(3)-structure on a five-dimensional Riemannian manifold $(M, g)$ is a structure defined by means of a smooth covariant completely symmetric and trace free tensor field $T$ of order 3 which satisfies on each chart $U(x^1, \ldots, x^5) \subset M$ the following condition:

$$T_{ijm} T_{kl}^m + T_{imk} T_{jl}^m + T_{mjk} T_{il}^m = g_{ij} g_{kl} + g_{ik} g_{jl} + g_{jk} g_{il}.$$  \hspace{1cm} (2.8)

For the tensor, $T \in C^\infty S^3_0 M$ the identities (2.8) after contraction with $g_{ij}$ and $T_{ij}^l$, respectively, imply

$$T_{kij} T_{lj}^k = \frac{7}{2} g_{ij}.$$ \hspace{1cm} (2.9)

$$T_{il}^m T_{jn}^n T_{km}^l = -\frac{3}{4} T_{ijk}.$$ \hspace{1cm} (2.10)

A five-dimensional Riemannian manifold $(M, g)$ with an irreducible SO(3)-structure was obtained as $(M, g, T)$. Further, authors in [10] define a characteristic $\nabla$ connection of $(M, g, T)$ with skew-symmetric torsion tensor such that $\nabla g = 0$; $\nabla T = 0$ and prove that an irreducible SO(3)-structure $(M, g, T)$ admits a characteristic connection $\nabla$ if it is nearly integrable, i.e., tensor field $T$ satisfies $\delta^* T = 0$. 

\[ Springer \]
connections

Using (3.2) and (3.3), we obtain the following identities:

\[
R^i_{jkl} = R^i_{jkl} + \frac{\alpha}{2}(\nabla_i T^j_{lk} - \nabla_k T^j_{li}) + \frac{\alpha^2}{4}(T^m T^i_{ml} - T^m T^i_{ml}), \quad (3.1)
\]

where \( R(X_k, X_l)X_j = R^i_{jkl}X_i \) and \( R(X_k, X_l)X_j = R^i_{jkl}X_i \) for \( X_k = \partial/\partial x^k \). Contracting with respect to the indices \( k \) and \( i \) in (3.1), we obtain the equalities

\[
\tilde{R}_{jl} = R_{jl} + \frac{\alpha}{2}(\nabla_i T^j_k - \nabla_k T^j_i) + \frac{\alpha^2}{4}(T^m T^i_{mj} - T^m T^i_{mj}), \quad (3.2)
\]

which connect the Ricci curvature tensor \( R^\alpha \) of the \( \alpha \)-connection \( \tilde{\nabla} \) and the Ricci curvature tensor \( Ric \) of the Levi-Civita connection \( \nabla \). In addition, we can rewrite the equalities in the form

\[
\tilde{-R}^\alpha = R_{jl} - \frac{\alpha}{2}(\nabla_i T^j_k - \nabla_k T^j_i) + \frac{\alpha^2}{4}(T^m T^i_{mj} - T^m T^i_{mj}). \quad (3.3)
\]

Equality (3.3) connects the Ricci curvature tensor \( Ric \) of the Levi-Civita connection \( \nabla \) and the Ricci tensor \( \tilde{Ric} \) of the connection \( \tilde{\nabla} \) which is a conjugate to \( \tilde{\nabla} \) relative to the metric \( g \). Using (3.2) and (3.3), we obtain the following identities:

\[
\tilde{R}^\alpha_{jl} - R^\alpha_{jl} = \alpha \left( \nabla_i T^j_k - \nabla_k T^j_i \right). \quad (3.4)
\]

Now we recall the definition of the Tchebychev form \( \tau \) [48, p. 58]. Namely, for conjugate connections \( \{ \nabla, g, \tilde{\nabla} \} \) the Tchebychev form \( \tau \) has local components \( T_k = g^{ij} T^k_{ij} \) for \( T_{ijk} = \Gamma^\alpha_{ijk} - \tilde{\Gamma}^\alpha_{ijk} \). Therefore, if \( \tilde{R}_{jl} = -R_{jl} \) then from (3.4) we have \( \nabla_j T_i = \nabla_i T_j \) and hence the \( \tilde{\nabla} \) is equiaffine. This shows that the Tchebychev form \( \tau \) is locally an exact 1-form \( df \) for some smooth scalar \( \text{potential function} \ f \) for \( \tau \).

On the other hand, the Christoffel symbols \( \Gamma^\alpha_{ij} \) of an equiaffine connection \( \tilde{\nabla} \) satisfy the equalities [48, pp. 57–59]

\[
\Gamma^\alpha_{kj} \equiv \partial_j \left( \ln \omega_{12...n} \right).
\]

But this time we know that the Christoffel symbols \( \Gamma^\alpha_{kj} \) of the Levi-Civita connection \( \nabla \) satisfy the following condition: \( \Gamma^\alpha_{kj} = \partial_j (\ln \sqrt{\det g}) \), where \( \omega = \sqrt{\det g} \ dx^1 \wedge \cdots \wedge dx^n \) is the Riemannian volume form of \( g \). Then using the above conditions for \( \Gamma^\alpha_{kj} \) and \( \Gamma^\alpha_{kj} \), we can transform equalities (2.1) into the differential equations

\[
\tilde{\partial}_k \ln \omega_{12...n} = \partial_k \ln \sqrt{\det g} + \frac{\alpha}{2} \partial_k f,
\]

where \( f \) is a potential function for \( \tau \). From this we can see that \( \alpha \omega_{12...n} = \sqrt{\det \tilde{g}} \), where \( \tilde{g} = \exp(C - \alpha f/n)g \) for an arbitrary constant \( C \).

Next, if we suppose that the Tchebychev form \( \tau = T_k dx^k \) is locally an exact 1-form, then from (3.2) and (3.3) we conclude that the \( \alpha \)-connection \( \tilde{\nabla} \) is equiaffine and its conjugate
connection $\nabla^\alpha$ is equiaffine, too. Next, multiplying by $g^{jl}$ on both sides of (3.2) yields
\[ \tilde{s} = s + \frac{\alpha^2}{4}(\|\tau\|^2 - \|T\|^2), \]
where $s = g^{jk}R_{jk}$ is the scalar curvature of $g$ and $\tilde{s} = g^{jk}\alpha^k$ is the scalar $\alpha$-curvature of $\nabla$. Hence, we have $\tilde{s} = s$ [48, p. 60] and if the Tchebychev form $\tau = 0$ then $\tilde{s} \leq s$. Finally, from the above considerations, we have the following.

**Theorem 3.1** Let $(M, g, T)$ be an $n$-dimensional statistical manifold. Assume that $\tau = \text{trace } T$ is the Tchebychev form of conjugate connections $\{\nabla, g, \nabla\}$, where $\nabla$ is a connection of one-parameter family of affine $\alpha$-connections on $M$. Then the following propositions hold.

1. $\tilde{-\alpha}s = \tilde{s} = s + \frac{\alpha^2}{4}(\|\tau\|^2 - \|T\|^2)$, where $\tilde{s}$ and $\tilde{-\alpha}s$ are scalar $\alpha$-curvature and ($-\alpha$)-curvature of $\nabla$ and $\nabla$, respectively.
2. If the Tchebychev form $\tau$ is locally an exact 1-form, then $\nabla$ and $\nabla$ are equiaffine for an arbitrary $\alpha$.
3. If $\text{Ric} = -\alpha\text{Ric}$, then the Tchebychev form $\tau$ is locally an exact 1-form with potential function $f$, $\bar{\nabla}$ is an equiaffine connection and the volume form $\tilde{\alpha}\omega = \sqrt{\text{det}g}$ for the Riemannian metric $\tilde{g} = \exp\left(\frac{C+\alpha f}{n}\right)g$, where $C$ is an arbitrary constant, which is conformal to $g$.

3.2 A conjugate symmetric statistical manifold

In this section, we consider a conjugate symmetric statistical manifold $(M, g, T)$. This means that $F = \nabla T$ is a covariant completely symmetric tensor field of order 4, i.e.,
\[ \nabla_k T_{lij} = \nabla_l T_{kij}, \tag{3.5} \]
where $T_{kij}$ are local components of $T$.

First, for a conjugate symmetric statistical manifold $(M, g, T)$, Theorem 3.1 implies the following corollary.

**Corollary 3.2** Each linear connection $\nabla^\alpha$ of the one-parameter family of $\alpha$-connections on an $n$-dimensional conjugate symmetric statistical manifold $(M, g, T)$ is an equiaffine connection such that the volume form $\tilde{\alpha}\omega$ associated with $\nabla$ is defined by the equality $\tilde{\alpha}\omega = \sqrt{\text{det}\tilde{g}}$ for $\tilde{g} = \exp\left(\frac{C+\alpha f}{n}\right)g$, where $C$ is an arbitrary constant.

Now let us consider a projectively flat equiaffine connection $\nabla$ for a fixed number $\alpha$ on a conjugate symmetric statistical manifold $(M, g, T)$. We recall that a torsion-free affine connection $\nabla$ on manifold $M$ is called projectively flat if there exists a flat affine connection $\nabla'$ on $M$ such that the pregeodesics (i.e., geodesics as curves without a distinguished parametrization) of $\nabla$ and $\nabla'$ coincide [35, pp. 138–142], [44, p. 386].

The $\alpha$-curvature tensor $\tilde{\alpha}R$ of a projectively flat equiaffine connection $\nabla$ has the form [35, pp. 138–142], [36, p. 169]
\[ \tilde{\alpha}R_{jkl} = \frac{1}{n - 1} \left( \delta_i^{\alpha} \tilde{\alpha}R_{jil} - \delta_i^{\alpha} \tilde{\alpha}R_{jki} \right). \tag{3.6} \]
Therefore, using identities (2.6) and the definition of the Ricci tensor $\alpha \Ric$, from (3.6) we deduce $\alpha R_{jk} = \frac{1}{n} \alpha s g_{jk}$. On the other hand, the Ricci tensor $\alpha \Ric$ is a Codazzi tensor with respect to a projectively flat equiaffine connection $\alpha \nabla$ [36, p. 169], i.e.,

$$\alpha \nabla_k R_{jl} = \alpha \nabla_j R_{lk}. \quad (3.7)$$

Then substituting $\alpha R_{jk} = \frac{1}{n} \alpha s g_{jk}$ into (3.7) we obtain $\alpha s = \text{const}$. Thus we have

$$\alpha R_{ijkl} = \frac{1}{n(n-1)} \alpha s (\delta_{ik} g_{jl} - \delta_{il} g_{jk}).$$

In this case, we say that $(M, g, T)$ is a statistical manifold of constant $\alpha$-curvature. For the case $\alpha = 1$ see [24,29,30].

It is well known that any Codazzi tensor $B$ of order 2 on a manifold $M$ with a projectively flat equiaffine connection $\alpha \nabla$ has local components [36, p. 169]

$$\alpha B_{ij} = \alpha \nabla_i \alpha \nabla_j f + \frac{f}{n-1} \alpha R_{ij}$$

for some smooth function $f$. On the other hand, we know that the metric tensor $g$ is a Codazzi tensor with respect to an arbitrary $\alpha$-connection $\alpha \nabla$. Thus, we have

$$g_{ij} = \nabla_i \alpha \nabla_j f + \frac{f}{n(n-1)} \alpha s \cdot g_{ij}$$

for $s = \text{const}$. It follows that $g_{ij} = \frac{n(n-1)}{n(n-1)-s} \alpha \nabla_i \alpha \nabla_j f$. We have now proved the following theorem.

**Theorem 3.3** Let $(M, g, T)$ be a conjugate symmetric statistical manifold. If there exists a projectively flat $\alpha$-connection $\alpha \nabla$ for some number $\alpha$, then $(M, g, T)$ is a statistical manifold of constant $\alpha$-curvature and its metric tensor $g$ has local components $g_{ij} = \frac{n(n-1)}{n(n-1)-s} \alpha \nabla_i \alpha \nabla_j f$ for some smooth function $f$ such that $(n(n-1)-s \cdot f) \nabla_i \alpha \nabla_j f$ is a positive definite symmetric 2-form on $TM$.

### 3.3 Compact statistical manifolds

First of all, we define a self-adjoint linear algebraic operator $\mathfrak{R} : S^2 M \to S^2 M$ acting on symmetric tensor fields of order 2 by the following identities:

$$\mathfrak{R}(B)_{ij} := R_{kij} B^{ij}.$$ 

This operator is called the curvature operator or curvature operator of second kind [9, pp. 51–52], [26,38,40].

From the above definition, the curvature operator $\mathfrak{R}$ is a self-adjoint linear operator and hence all its eigenvalues are real numbers. This curvature operator is called negative or positive definite operator if its eigenvalues are negative or positive numbers, respectively.

It is well known [9, p. 45] that the vector space $S^2 M$ of covariant symmetric tensor fields of order 2 is irreducible with respect to the action of the orthogonal group $O(n)$ at each point $x \in M$. We denote by $S^2_0 M$ the vector space of traceless symmetric tensor fields of order 2 over $(M, g)$. Then the following pointwise $O(n)$-irreducible decomposition holds:
Lemma 3.4 Let \((M, g, T)\) be a compact oriented statistical manifold. Then the following integral formula holds:

\[
\langle Ric(T), T \rangle + 2\langle \mathcal{H}(T), T \rangle + \frac{1}{12}(\delta^* T, \delta^* T) - \frac{1}{3}\langle F, F \rangle - \langle \delta T, \delta T \rangle = 0. \tag{3.8}
\]

Proof Let \(X\) be a vector field with local components \(T^l_{jk} \partial_j T^{ijk} - T^i_{jk} \partial_i T^{ijk} + T^i_{jk} \partial_j T^{ijl} - T^i_{jk} \partial_i T^{ijl}\). For this vector field we have

\[
\text{div} X = \nabla_i T^l_{jk} \nabla_j T^{ijk} + T^l_{jk} (\nabla_i \nabla_j T^{ijk}) - \nabla_j T^l_{jk} \nabla_j T^{ijl} - T^i_{jk} (\nabla_i T^{ijk})
\]

\[
= R_{ij} T^{ikl} T^j_{kl} + 2R_{ijkl} T^{lm} T^k_m + \nabla_i T^{jkl} \nabla_j T^i_{kl} - \nabla_j T^{ikl} \nabla_m T^m_{kl},
\]

\[
= g(Ric(T), T) + 2g(\mathcal{H}(T), T) + \frac{1}{12}g(\delta^* T, \delta^* T) - \frac{1}{3}g(F, F) - g(\delta T, \delta T), \tag{3.9}
\]

where we used the Ricci identities [28, p. 145]

\[
\nabla_i \nabla_j T^{ikl} - \nabla_j \nabla_i T^{ikl} = T^{mkl} R^i_{mij} + T^{iml} R^k_{mij} + T^{ikm} R^l_{mij}, \tag{10.10}
\]

and defined \(g(Ric(T), T)\) and \(g(\mathcal{H}(T), T)\) by the equalities \(g(Ric(T), T) = R_{ij} T^{ikl} T_{jl}^i\) and \(g(\mathcal{H}(T), T) = (R_{ijkl} T^{ilm}) T^k_m\), respectively.

If we suppose that \(M\) is a compact manifold, then applying Green’s theorem [28, p. 281] to (3.9) we obtain the integral formula (3.8). Here, we note that, to apply Green’s theorem, it is necessary to assume \(M\) as orientable. If \(M\) is not orientable, then we only need to consider an orientable double covering. The proof is complete. \(\square\)

Next, using Lemma 3.4 we prove the following “vanishing theorem”.

Theorem 3.5 Let \((M, g, T)\) be a compact oriented statistical manifold. If the curvature operator \(\mathcal{H}\) is positive definite on \(S^2_0 M\), then \((M, g, T)\) cannot be a conjugate symmetric statistical manifold with vanishing Tchebychev form. In particular, an \(n\)-dimensional \((n \geq 3)\) Euclidean sphere \((S^n, g_0)\) cannot be a conjugate symmetric statistical manifold with vanishing Tchebychev form.

Proof Let \((M, g, T)\) be a conjugate symmetric statistical manifold then follows the equality \(\delta^* T = 4F\), where \(F\) is given by (2.5). On the other hand, from (3.5), we obtain \(\delta T = -\nabla \tau\) for the Tchebychev form \(\tau\) of \((M, g, T)\). Therefore, if \((M, g, T)\) is a compact conjugate symmetric statistical manifold with vanishing Tchebychev form \(T\) then \(T \in C^\infty S^2_0 M\) and we can rewrite (3.8) in the form

\[
\langle Ric(T), T \rangle + 2\langle \mathcal{H}(T), T \rangle = -\langle F, F \rangle \leq 0. \tag{3.11}
\]
If we suppose that the curvature operator $\mathfrak{R}$ is positive definite on $S_0^2 M$, i.e.,

$$R_{ijkl} B^{lk} B^{il} \geq \lambda B_{ij} B^{ij} > 0 \quad (3.12)$$

for an arbitrary symmetric tensor field $B = (B_{ij})$ which belongs to $S_0^2 M$ and some positive number $\lambda$, then

$$g(\mathfrak{R}(T), T) \geq \lambda \|T\|^2 > 0.$$ 

Denoting by $Y \circ X$ the symmetric tensor $X \circ Y : = 1/2(X \otimes Y + Y \otimes X)$ for any local orthogonal unit vector fields $X$ and $Y$, we can represent (3.10) as $R_{ijkl} X^i Y^j X^k Y^l \geq \lambda > 0$. This means that the sectional curvature $sec(X \wedge Y) \geq \lambda > 0$ if the curvature operator $\mathfrak{R} \geq \lambda > 0$ on $S_0^2 M$ [38, 40].

Next for the $n - 1$ local unit vector fields $Y_1, \ldots, Y_{n-1}$, orthogonal to $X$ and to each other, we have $sec(X \wedge Y_a) = R_{ijkl} Y^i_a X^j X^k Y^l_a \geq \lambda > 0$ for all $a = 1, \ldots, n - 1$, and hence $Ric (X, X) \geq (n - 1)\lambda > 0$ [40]. Then we claim that

$$g(\text{Ric}(T), T) \geq (n - 1)\lambda \|T\|^2 > 0.$$ 

So the left-hand side of (3.11) is positive, and this contradicts the inequality (3.11).

In particular, let $(M, g)$ be a Riemannian manifold of constant sectional curvature then the Riemann curvature tensor has the form $R_{ijkl} = K (g_{ij} g_{lk} - g_{ik} g_{lj})$. In this case (3.8) can be rewritten as

$$(n - 3) K \langle T, T \rangle = -\langle F, F \rangle \leq 0. \quad (3.13)$$

For $n > 3$, $K > 0$ and $T \neq 0$ the left-hand side of (3.13) is positive that contradicts the inequality (3.13). In particular, for $n = 3$ from (3.13) we obtain $F = 0$. In this case, the Ricci identities (3.10) can be rewritten in the form $(n + 1)K T^k_l = 0$ for $K > 0$ and $T \neq 0$, and this is a contradiction. This completes the proof of Theorem 3.5. \hfill $\square$

We shall now consider a harmonic statistical manifold $(M, g, T)$ and show that the following theorem is true.

**Theorem 3.6** Let $(M, g, T)$ be a compact oriented statistical manifold. If the curvature operator $\mathfrak{R}$ is negative definite, then $(M, g, T)$ cannot be a harmonic statistical manifold. In particular, an $n$-dimensional $(n \geq 3)$ compact hyperbolic manifold $(\mathbb{H}^n, g_0)$ cannot be a harmonic statistical manifold.

**Proof** Let $(M, g, T)$ be a compact harmonic statistical manifold. In this case the formula (3.12) shows that

$$\langle Ric(T), T \rangle + 2 \langle \mathfrak{R}(T), T \rangle = \frac{11}{12} \langle \delta^* T, \delta^* T \rangle + \frac{1}{3} \langle F, F \rangle \geq 0, \quad (3.14)$$

because $\langle \Delta_{\text{sym}} T, T \rangle = \langle \delta T, \delta T \rangle - \langle \delta^* T, \delta^* T \rangle$. If we suppose that the curvature operator $\mathfrak{R}$ is negative definite, i.e.,

$$R_{ijkl} B^{lk} B^{il} \leq -\mu \cdot B_{ij} B^{ij} < 0$$

for an arbitrary symmetric tensor field $B = (B_{ij})$ and some positive number $\mu$, then $g(\mathfrak{R}(T), T) \leq -\mu \cdot \|T\|^2 < 0$ and hence $g(\text{Ric}(T), T) \leq -(n - 1)\mu \cdot \|T\|^2 < 0$ [40]. In this case the left-hand side of (3.14) can be rewritten in the form

$$\langle Ric(T), T \rangle + 2 \langle \mathfrak{R}(T), T \rangle = -(n + 1)\mu \langle T, T \rangle \leq 0.$$ 

So the left-hand side of (3.14) is negative that contradicts the inequality (3.14).
In particular, let \((M, g)\) be a Riemannian manifold of constant sectional curvature \([18, \text{p. 84}]\) then \(R_{ij} = K (g_{ij} g_{lk} - g_{ik} g_{lj})\). In this case (3.14) can be rewritten in the form

\[
K ((n - 3) (T, T) + 2 (\tau, \tau)) = \frac{11}{12} \langle \delta^* T, \delta^* T \rangle + \frac{1}{3} (F, F) \geq 0. \tag{3.15}
\]

For \(n > 3, K > 0\) and \(T \neq 0\) the left-hand side of (3.15) is negative, and this contradicts the inequality (3.15). In particular, for \(n = 3\) from (3.15) we obtain \(F = 0\) and \(\tau = 0\). In this case the Ricci identities (3.8) can be rewritten in the form \((n + 1) K T^l_{ij} = 0\) for \(K > 0\) and \(T \neq 0\). As a result, we have a contradiction. This completes the proof of Theorem 3.5. \(\Box\)

4 A five-dimensional Riemannian manifold with an irreducible SO(3)-structure

4.1 A conjugate symmetric statistical manifold with an irreducible SO(3)-structure

Let \((M, g, T)\) be a five-dimensional statistical manifold with an irreducible SO(3)-structure, this means that the skewness tensor \(T\) of \((M, g, T)\) defines an irreducible SO(3)-structure on \((M, g)\). This shows that the skewness tensor \(T\) of \((M, g, T)\) belongs to \(S^3_0 M\) (i.e., the Tchebychev form \(\tau\) of \((M, g, T)\) is zero).

Consider a one-parameter group of affine torsion-free connections \(\mathbf{\alpha} \nabla\) on \((M, g)\) by Eq. (2.1). Then using (2.8) and (2.9) we get the equations

\[
\nabla_{l}^{(\alpha)} T_{ijk} = \nabla_{l} T_{ijk} + \frac{\alpha}{2} (T_{ijm} T_{m^l} + T_{imk} T_{m^j} + T_{mj} T_{m^i}^l) = \nabla_{l} T_{ijk} + \frac{\alpha}{2} (g_{ij} g_{kl} + g_{ik} g_{jl} + g_{jk} g_{il}). \tag{4.1}
\]

Then from Eq. (4.1), we obtain the following equality:

\[
\| \nabla T \|^2 = \| F \|^2 + \frac{105}{4} \alpha^2, \tag{4.2}
\]

where \(F\) is given by (2.5), \(\| \nabla T \|^2 = g(\nabla T, \nabla T)\) and \(\| F \|^2 = g(F, F)\). Then the tensor \(T\) of an irreducible SO(3)-structure satisfies the following inequalities:

\[
\| \nabla T \|^2 \geq \| F \|^2; \quad \| \nabla T \|^2 \geq \frac{105}{4} \alpha^2 > 0. \tag{4.3}
\]

On the other hand, from (3.2) we obtain the following identities:

\[
\nabla_{\alpha} R_{jk} = R_{jk} - \frac{\alpha}{2} \nabla_{m} T_{m^l} - \frac{7 \alpha^2}{4} g_{jk}. \tag{4.4}
\]

In this case, \(\nabla_{\alpha} R_{jk} = R_{kj}\) and hence \(\nabla\) is an equiaffine connection for an arbitrary \(\alpha\). Next from (4.4), we obtain \(s = s - 35/4 \cdot \alpha^2\) and hence \(s \leq s\) for the scalar \(\alpha\)-curvature \(s\) of \(\nabla\) and scalar curvature \(s\) of \(g\). The next theorem is an immediate consequence of Theorem 3.1 and its proof.

**Theorem 4.1** Let a five-dimensional Riemannian manifold \((M, g)\) with an irreducible SO(3)-structure be a statistical manifold \((M, g, T)\), then the one-parameter family of affine \(\alpha\)-connections \(\nabla^{\alpha}\) consisting of equiaffine connections and their volume forms \(\omega^{\alpha} = \sqrt{\det(\exp C \cdot g)}\) do not depend on the parameter \(\alpha\). Moreover, the scalar \(\alpha\)-curvature \(\overset{\alpha}{\nabla}\) has the form \(\overset{\alpha}{s} = s - \frac{35}{4} \cdot \alpha^2\). \(\Box\)

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We consider a conjugate symmetric statistical manifold \((M, g, T)\) such that the skewness tensor field \(T\) defines an irreducible SO(3)-structure on \((M, g)\). In this case, the skewness tensor field \(T\) satisfies the following equations:

\[
\nabla_k T_{ljk} = \nabla_l T_{kij}; \quad \nabla^k T_{kij} = 0,
\]

where \(\nabla^k = g^{kl} \nabla_l\). In view of what we have told above, we can formulate the following theorem.

**Theorem 4.2** Let \((M, g, T)\) be a five-dimensional statistical manifold with an irreducible SO(3)-structure and \(\alpha\) a one-parameter family of affine \(\alpha\)-connections on \(M\). Then \((M, g, T)\) is conjugate symmetric if and only if \(\alpha \nabla T\) is a covariant completely symmetric tensor field of order 4. In particular, if \((M, g, T)\) is a statistical manifold of constant \(\alpha\)-curvature, then it is a conjugate symmetric statistical manifold with Einstein metric \(g\).

**Proof** From (3.5), we obtain the following identities:

\[
\alpha \nabla_l T_{ijk} - \alpha \nabla_i T_{ljk} = \nabla_l T_{ijk} - \nabla_i T_{ljk}.
\]

In turn, from (4.6) we conclude that a five-dimensional Riemannian manifold \((M, g)\) with an irreducible SO(3)-structure is a conjugate symmetric statistical manifold if and only if the condition \(\nabla^\alpha T\) is a completely symmetric tensor field of order 4.

In addition, we suppose that \((M, g, T)\) is a statistical manifold of constant \(\alpha\)-curvature for some \(\alpha\) then from (2.3), (2.8) and the Ricci identities [18, p. 30]

\[
\alpha \nabla_k \alpha \nabla_l g_{ij} - \alpha \nabla_l \alpha \nabla_k g_{ij} = -g_{lm} \alpha \nabla_l \alpha \nabla_m T_{jkl} + \|F\|^2 = \frac{7}{2} s + \frac{2}{7} \alpha^2 \cdot g_{jk}.
\]

we obtain (4.5). In addition, from (4.4) we have \(R_{jk} = \left(\frac{1}{5} \frac{\alpha^2}{2} \right) \cdot g_{jk}\). This means that \(g\) is an Einstein metric. The proof is complete.

Next we shall prove two “vanishing theorems” for conjugate symmetric statistical manifolds.

**Theorem 4.3** Let \((M, g, T)\) be a five-dimensional statistical manifold such that \(T\) be a structure tensor of an irreducible SO(3)-structure on \((M, g)\). If the curvature operator \(\Re (M, g)\) is positive definite on \(S_0^2 M\), then \((M, g, T)\) cannot be a (locally) conjugate symmetric. In particular, a five-dimensional Euclidian sphere \((S^5, g_0)\) cannot be a (locally) conjugate symmetric statistical manifold with an irreducible SO(3)-structure.

**Proof** Let \((M, g, T)\) be a conjugate symmetric statistical manifold. Then from the identity (2.8) by straightforward computation we get

\[
0 = \nabla_i \nabla_j (T^{ikl} T_{kjl}) = (\nabla_i \nabla_j T^{ikl}) T_{kjl} + (\nabla_j T^{ikl}) (\nabla_i T_{kjl})
\]

\[
= (T^{mkj} R_{mjkl}^i + T^{imj} R_{mkjl}^k + T^{ikm} R_{mklj}^i) T_{kjl} + \|F\|^2 = \frac{7}{2} s + 2 g(\Re (T), T) + \|F\|^2,
\]

where \(F\) is given by (2.5). If we suppose that the curvature operator \(\Re\) is positive definite on \(S_0^2 M\), then the scalar curvature \(s\) of \((M, g)\) is positive, too [40]. Moreover, in this case, we have \(g(\Re (T), T) > 0\). Thus, we have the inequality \(\frac{7}{2} s + 2 g(\Re (T), T) + \|F\|^2 > 0\) which contradicts to (4.8). The proof is complete.
The identities (3.1) for a conjugate symmetric statistical manifold \((M, g, T)\) have the form
\[
R_{ijkl} = \alpha R_{ijkl} + \frac{\alpha^2}{4} (T_{mil}T_{mj} - T_{mji}T_{ik}),
\]
where \(R_{ijkl} := g_{im} \alpha R_{mjk}^l\). From (2.6) we conclude that symmetry properties of \(R_{ijkl}\) coincide with the symmetry properties of the Riemann curvature tensor \(R_{ijkl}\). Therefore, we can define the \(\alpha\)-curvature operator \(\alpha \mathbb{R}\) of an \(\alpha\)-connection \(\alpha \nabla\) as an analog of the curvature operator \(\mathbb{R}\) of \((M, g)\). Then using (2.8), (2.9) and (2.10) from (4.9) we get
\[
g(\mathbb{R}(T), T) = g(\mathbb{R}(\mathfrak{H}(T), T) + \frac{505}{16} \alpha^2.\]
If we suppose that there is an \(\alpha\)-connection \(\alpha \nabla\) such that its curvature operator \(\alpha \mathbb{R}\) is positive semi-definite then \(g(\mathbb{R}(T), T) \geq \frac{505}{16} \alpha^2 > 0\). Thus, as a consequence of Theorem 4.3, we have the following corollary.

**Corollary 4.4** Let \((M, g, T)\) be a five-dimensional Riemannian manifold \((M, g)\) with an irreducible SO(3)-structure. If there exists an \(\alpha\)-connection \(\alpha \nabla\) such that its curvature operator \(\alpha \mathbb{R}\) is positive semi-definite then \((M, g, T)\) is not conjugate symmetric.

Finally, we present topological properties of a five-dimensional conjugate symmetric statistical manifold with an irreducible SO(3)-structure. Namely, we have the following Theorem.

**Theorem 4.5** Let \((M, g, T)\) be a conjugate symmetric statistical manifold such that \((M, g)\) is a five-dimensional compact, conformally flat Riemannian manifold and \(T\) is a structure tensor of an irreducible SO(3)-structure on \((M, g)\). Then the fundamental group \(\pi_1(M)\) has exponential growth unless \(F \equiv 0\). In particular, if \(\pi_1(M)\) is almost solvable, then \((M, g)\) is a flat manifold. Moreover, if \(F\) does not vanish anywhere on \((M, g)\) then \(\pi_1(M)\) is a non-amenable group.

**Proof** The Riemannian curvature tensor \(R\) of a five-dimensional conformal flat Riemannian manifold \((M, g)\) has the form [18, p. 92]
\[
R_{ijkl} = \frac{1}{3} (g_{ik}R_{lj} - g_{il}R_{jk} + g_{lj}R_{ik} - g_{jk}R_{il}) + \frac{s}{12} (g_{jl}g_{ik} - g_{jk}g_{li}).
\]
In this case from (4.8) we obtain \(s = -\frac{4}{3} \|F\|^2 \leq 0\), where \(F\) is given by (2.5).

We know three things: First, if \((M, g)\) is compact, conformally flat with non-positive scalar curvature \(s\), then the fundamental group \(\pi_1(M)\) has exponential growth unless \(s \equiv 0\) [31]. This proposition gives the first result. Second, if the fundamental group \(\pi_1(M)\) of a compact conformal flat Riemannian manifold \((M, g)\) with non-positive scalar curvature is almost solvable, then \((M, g)\) is a flat manifold [31]. This proposition gives the second result. Third, the fundamental group \(\pi_1(M)\) of a compact, locally conformally flat Riemannian manifold \((M, g)\) with negative scalar curvature is a non-amenable group [46]. That gives the third result of the theorem.

\(\Box\)

4.2 Manifolds with nearly integrable SO(3)-structures

In this section, we consider a five-dimensional Riemannian manifold \((M, g)\) with a nearly integrable SO(3)-structure as an example of a harmonic statistical manifold. In this case, we have the following theorem.

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Theorem 4.6 A five-dimensional Riemannian manifold with a nearly integrable SO(3)-structure is a harmonic statistical manifold such that its skewness tensor defines the first integral of order 3 of the equations of geodesics.

Proof Let $T$ be the tensor of a nearly integrable SO(3)-structure on a five-dimensional Riemannian manifold $(M, g)$. If the equation $\delta^* T = 0$ holds, then necessarily $\delta T = 0$ as is readily seen, and so $\Delta_{\text{sym}} T := \delta^* \delta T - \delta \delta^* T = 0$, where $\Delta_{\text{sym}}$ is a Laplacian on symmetric covariant tensor fields [9, pp. 52–54], [45]. In this case, the tensor field $T$ is harmonic [45] and hence, $(M, g, T)$ is a harmonic statistical manifold.

On the other hand, a 1-dimensional immersed submanifold $\gamma$ of $(M, g)$ is said to be a geodesic if there exists a parameterization $\gamma : x^i = x^i(t)$ for $t \in I \subset \mathbb{R}$ satisfying $\nabla_\gamma \dot{x} = 0$. If each solution $x^i = x^i(t)$ of the equations $\nabla_\gamma \dot{x} = 0$ of geodesics satisfies the condition $B(\dot{x}, \ldots, \dot{x}) = \text{const}$ for smooth covariant completely symmetric tensor field $B$ of order $p$ and $\dot{x} = \frac{dx^k}{dt} \cdot X_k$, then the equations $\nabla_\gamma \dot{x} = 0$ are said to admit the first integral of the $p$-th order. The equation $\delta^* B = 0$ serves as a necessary and sufficient condition for this [18, pp. 128–129].

It is easy to check that a five-dimensional statistical manifold $(M, g, T)$ whose skewness tensor $T$ is a structure tensor of a nearly integrable irreducible SO(3)-structure on $(M, g)$ admits the first integral of order 3 of geodesics in the form $T(\dot{x}, \dot{x}, \dot{x}) = \text{const}$. Thus, the theorem is valid. □

Remark For a harmonic symmetric tensor, it was said in [45] that it does not have any geometrical interpretation. It was also pointed out that symmetric tensors, whose covariant derivative vanishes, are clearly harmonic, the metric tensor $g$ being the most important example. These tensors are trivial examples of a harmonic symmetric tensor. Now we can say that the tensor $T$ of a nearly integrable SO(3)-structure on a five-dimensional connected Riemannian manifold $(M, g)$ is a non-trivial example of a harmonic symmetric tensor.

Next, we prove a “vanishing theorems” for nearly integrable SO(3)-structures.

Theorem 4.7 Let the triplet $(M, g, T)$ be a statistical manifold, where $(M, g)$ be a five-dimensional Riemannian manifold and $T$ be a structure tensor of an irreducible SO(3)-structure on $(M, g)$. If the curvature operator $\mathfrak{H}$ of $(M, g)$ is negative definite, then SO(3)-structure cannot be nearly integrable. In particular, there is no a nearly integrable irreducible SO(3)-structure on a five-dimensional hyperbolic manifold $(H^5, g_0)$.

Proof We consider a five-dimensional statistical manifold $(M, g, T)$ such that its skewness tensor $T$ defines a nearly integrable irreducible SO(3)-structure on $(M, g)$. In this case, from the identity (2.8) by straightforward computation, we get

$$0 = \nabla_i \nabla_j (T^{ikl} T_{kl}^j) = (\nabla_i \nabla_j T^{ikl}) T_{kl}^j + (\nabla_i T^{ikl}) (\nabla_j T_{kl}^j)$$

$$= (T^{mkl} R_{mij} + T^{ilm} R_{mij}^k + T^{ikm} R_{mij}^l) T_{kl}^j - \frac{1}{3} \| F \|^2$$

$$= \frac{7}{2} s + 2 g(\mathfrak{H}(T), T) - \frac{1}{3} \| F \|^2. \quad (4.10)$$

If we suppose that the curvature operator of second kind $\mathfrak{H}$ is negative definite, then $g(\mathfrak{H}(T), T) < 0$ and the scalar curvature $s$ must be negative, too. The above conditions give rise to a contradiction with the formula (4.10).

In particular, if we consider a Riemannian manifold $(M, g)$ of constant curvature $K$, then from the formula (4.10) we can obtain $K \geq 0$. Hence there is no a nearly integrable
irreducible SO(3)-structure on a five-dimensional hyperbolic manifold \((H^5, g_0)\) [55]. This completes the proof of the theorem.

Finally, we consider a topological property of a five-dimensional conformally flat \((M, g)\) with a nearly integrable irreducible SO(3)-structure. Namely, we have the following corollary.

**Corollary 4.8** Let \((M, g)\) be a compact, locally conformally flat Riemannian manifold with a nearly integrable irreducible SO(3)-structure whose structure tensor \(T\) has a non-vanishing covariant derivative \(\nabla T\) on \(M\). Then the fundamental group \(\pi_1(M)\) is hyperbolic in the sense of Gromov.

**Proof** For a locally conform flat Riemannian manifold \((M, g)\) from (4.10), we obtain

\[
s = \frac{4}{63} \|F\|^2 \geq 0.
\]

As a consequence of main theorem [25], we have that the fundamental group \(\pi_1(M)\) of a compact, conformally flat manifold \((M, g)\) with positive scalar curvature is hyperbolic in the sense of Gromov. This proposition gives our result. In addition, we say that there are other applications to geometry and topology of locally conformally flat manifolds with positive scalar curvature, and hence to geometry and topology of locally conformally flat Riemannian manifold with a nearly integrable irreducible SO(3)-structure [25]. In particular, this is a “vanishing theorem” for cohomology groups of such manifolds.

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