LINEARIZATION OF SKEW-PERIODIC LOOPS AND S1-COCYCLES

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ABSTRACT. We discuss linearization of skew-periodic loops. We generalize the situation to linearization of non-commutative loops and S1-cocycles.

INTRODUCTION

Linearization is a phenomenon when a nonlinear system is made linear at the cost of adding extra space-dimensions but retaining essential equivalence to the old system. The best known example is that there is a homotopy

\[
\begin{bmatrix}
a_0 + a_1 z + \ldots + a_n z^n \\
1 \\
\vdots \\
a_{n-1} \\
a_n
\end{bmatrix}
\sim
\begin{bmatrix}
a_0 \\
a_1 \\
a_2 \\
\vdots \\
a_{n-1} \\
a_n
\end{bmatrix}
\begin{bmatrix}
-z \\
1 \\
-1 \\
\ddots \\
1 \\
1
\end{bmatrix}
\]

such that this homotopy is achieved via multiplication by invertible matrices. In particular, if \(a_0 + a_1 z + \ldots + a_n z^n\) was invertible then the homotopy will be through invertible matrices. In case of non-polynomial systems one has to use infinitely many dimensions and it may not be clear how to proceed even for convergence reasons.

In terms of loops the linearization problem is like to transform an invertible Laurent loop \(a(z) = \sum_{n=-\infty}^{\infty} a_n z^n\) into a linear loop \(\bar{a}(z) = \bar{a}_0 + \bar{a}_1 z\) by a homotopy. In [2] a sufficiently nice method was presented for the case just mentioned, which is essentially the case of the pointed complex loops.

In this paper we look a bit further, we show how to discuss the case of real loops, and some other ones. As it turns out the most general case is not necessarily the most enlightening, so we will start with specific examples, and then we proceed further to the more abstract setting. The most basic case is the case of what we call based loops, ie. maps from \(S^1\) into the unit group of an algebra. Having discussed this case we turn to various generalizations. One possible generalization is when we use non-commutative loops in the sense such that the loop variable \(z\) would not commute with the coefficients \(a_n\). Another generalization is about linearization of \(S^1\)-cocycles, as explained below, which allows to treat generalized loops which are like maps from a principal \(S^1\)-bundle into a unit algebra, or a regular \(S^1\)-action.

Another thing is that we concentrate on linearizing skew-periodic loops. Compared to the previous example this is like to try to linearize the loop \(a(z) = \sum_{n \in \mathbb{Z}+\frac{1}{2}} a_n z^n\) into the linear loop \(\bar{a}(z) = \bar{a}_{-\frac{1}{2}} z^{-1/2} + \bar{a}_{\frac{1}{2}} z^{1/2}\). This is a more

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natural thing to do. Our methods here are the essentially the same ones as in [2] but we feel worthwhile to present the modifications necessary.

Let us summarize here how we expect to execute linearizatizations: It is assumed to be done in two steps: 1. We add extra space-variables. This is supposed to be done by a trivial extension. 2. We linearize by a homotopy. Now, there are some requirements: a.) If a loop was already linear then linearization should essentially do nothing. b.) The construction should be sufficiently natural.

The case of based loops is somewhat misleading. Having a base point is like having an anchor which allows linearization of loops as such. But even there, what we really do is the linearization of the cocycle $\tilde{a}(z, w) = a(z)a(w)^{-1}$. In the more general case we should rather linearize cocycles obtained not necessarily from loops. This is in accordance to the geometric picture: when we apply the clutching construction for vector bundles then what we need is not necessarily a clutching function but a clutching cocycle in terms of the structure group. There is an other widely used way to represent cocycles: that is to take the corresponding element from the skew-loop Grassmannian space (cf. [1]). As it turns out, the corresponding Grassmannian element directly refers to the linearized cocycle. In particular, we do not lose information in the linearization process.

In Section 1 we consider some very basic matrix constructions associated to based loops. In order to represent loops we consider skew-complex and real Fourier bases. In Section 2 we define some auxiliary constructions, they are basically about how to write down some matrices naturally. In Section 3 we present some constructions which are more to the point: some deformation matrices are introduced. Using them, in Section 4 some very natural deformations of multiplication actions are introduced. Section 5 is a short digression about deformations of rapidly decreasing matrices in similar manner. In Section 6, at last, we present the linearizing homotopies for based loops. By that time we will have more than enough information to see why these linearizations work. Section 7 is about generalization to non-commutative loops. Section 8 is a more detailed discussion of transformation kernels on $S^1$ which are useful alternatives to matrices, and which may motivate some of our constructions in the next, last section: Section 9 is about linearization of $S^1$-cocycles. It sort of embodies the naturality of the basic linearization construction. As it yields a somewhat abstract setting it is useful to think through its possible geometric interpretations.

Regarding analysis: We will treat the smooth case, even if we make efforts to point into the direction of algebraic generalizations. Here “smooth” is used in the sense of infinitely many differentiable or rapidly decreasing (in case we consider a Fourier or matrix expansions). People with geometric applications on their mind may know what modifications are required in specific cases.

Some of the length of this paper is gained from that we try to show the inner workings of the constructions. In particular, most of our constructions come in three varieties: the left perturbation case ($L$), the “unitary” case ($\mathcal{O}$), and the right perturbation case ($R$). (The polynomial linearizing homotopy above would be of type “$L$.”) It would not be hard to give a much shorter exposition of the “unitary” case alone, but some generalizations fit more to the “$L$” or “$R$” cases.

In terms of proofs, some of them are tedious matrix computations. The author has decided not to include them here but to present those key points which give the frame of the computations making them straightforward.
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1. Matrices associated to based loops

1.1. In what follows \( \mathfrak{A} \) will be a locally convex algebra, \( 1 \in \mathfrak{A} \), although it is not hard to transcribe the constructions to the non-unital case. We can define the opposite algebra \( \mathfrak{A}^{\text{opp}} \) such that to any \( X \in \mathfrak{A} \) we associate the opposite element \( X^{\text{opp}} \in \mathfrak{A}^{\text{opp}} \). This operation is linear but opposite elements multiply in a contravariant manner: \( X^{\text{opp}} \cdot Y^{\text{opp}} = (YX)^{\text{opp}} \).

A more usual manifestation of this operation is when \( \mathfrak{A} \) contains a distinguished skew-involution \( g \) and there is an isomorphism \( \psi : \mathfrak{A}^{\text{opp}} \rightarrow \mathfrak{A} \) given such that \( \psi(g^{\text{opp}}) = -g \). Then the star operation \( a^* = \psi(a^{\text{opp}}) \) can be used, it yields a \( g \)-antilinear operation.

1.2. Notation. In order to simplify the notation, we make the following convention: In the associated constructions instead of \( "x \mapsto a(x)" \) we will often write simply \( "a" \). This assumes that the variable of \( a \) is \( x \), and we do this only if there is no danger of confusion.

Matrices will be considered as linear combinations of elementary matrices. General elementary matrices \( e_{a,m} = e_n \otimes e_m^\top \) are considered as tensor products of elementary column and row matrices. In what follows \( "\cdot\cdot\cdot" \) denotes a diagonal pattern of something and \( "\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\cdot\·\cdot\"\) denotes a diagonal pattern of 1’s. If \( S \) is an appropriate set of indices then we define the identity, flip, and shift matrices as

\[
1_S = \sum_{s \in S} e_{s,s}, \quad T_{-S,S} = \sum_{s \in S} e_{-s,s}, \quad Z_{n,S,S}^S = \sum_{k \in S} e_{n,k,k},
\]

respectively. If two lower indices are the same then only one is used. We define \( S^+ = \{s \in S : s > 0\} \), etc. We will mainly be interested in the cases \( S = \mathbb{Z} \) and \( S = \mathbb{Z}_i = \mathbb{Z} + \frac{1}{2} \). In case we do not fix which one we use we apply the notation \( Z_i, i \in \mathbb{Z}_i = \mathbb{Z}/2\mathbb{Z} \) where \( \mathbb{Z}_0 = \mathbb{Z} \) and \( \mathbb{Z}_1 = 2\mathbb{Z} \). The complementer set is \( Z_i^c = (\frac{1}{2}\mathbb{Z}) \setminus Z_i = Z_i + \frac{1}{2} \). When writing \( \mathbb{Z} \times \mathbb{Z} \) matrices we draw two lines, above and below the 0th row, when writing \( \mathbb{Z}_i \times \mathbb{Z}_i \) matrices we draw one line between the \( \frac{1}{2} \)th and \((-\frac{1}{2})\)th row; similarly with columns.

If \( \phi : \mathfrak{A} \rightarrow \mathfrak{A}' \) is a homomorphism and \( \mathbf{A} \) is a \( \mathfrak{A} \)-valued matrix then let simply \( \phi(\mathbf{A}) \) denote the matrix which is obtained from \( \mathbf{A} \) by replacing every entry \( a_{n,m} \) with its homomorphic image \( \phi(a_{n,m}) \). The generalization of taking opposites is given by taking \( \mathbf{A}^{\text{opp}} \), i.e. taking the transpose of \( \mathbf{A} \) and replacing every entry \( a_{m,n} \) by \( a_{n,m}^{\text{opp}} \).

1.3. Loops and Fourier expansions. Let \( \mathfrak{B}^0 = C_\mathbb{Z}(\mathfrak{A}) \) be the space of smooth periodic \( \mathfrak{A} \)-valued loops, i.e. of functions \( a : \mathbb{R} \rightarrow \mathfrak{A} \) such that \( a(x + 2\pi) = a(x) \). Similarly, let \( \mathfrak{B}^1 = C_\mathbb{Z}(\mathfrak{A}) \) be the space of smooth skew-periodic \( \mathfrak{A} \)-valued loops, i.e. of functions \( a : \mathbb{R} \rightarrow \mathfrak{A} \) such that \( a(x + 2\pi) = -a(x) \). Then \( \mathfrak{B} = \mathfrak{B}^0 \oplus \mathfrak{B}^1 \) is a naturally \( \mathbb{Z}_2 \)-graded algebra. Such loops allow Fourier expansions in various ways:

a.) Suppose that \( g \) is a skew-involution. Let us define \( z(x) = \cos x + g \sin x \) where \( x \in \mathbb{R} \). We use the notation \( z^n(x) = \cos nx + g \sin nx \). Then \( a \in \mathfrak{B}^n \) allows a Fourier expansion \( a(x) = \sum_{n \in \mathbb{Z}_i} z^n(x) a_{n,i} \). Using the correspondence \( e_n \leftrightarrow z^n(x) \) of bases we
can define the column matrix
\[ S_{Z_i}^g(a) = \sum_{n \in Z_i} a_n e_n. \]

Similarly, we can also define the row matrices \( S_{Z_i}^g(a) \): For a skew-involution \( g \) we will use the short notation \( [K]_g^+ = \frac{K+gKg}{2} \), \( [K]_g^- = \frac{K-gKg}{2} \), for the \( g \)-invariant and \( g \)-skew-invariant parts of \( K \), respectively. We see that \( [K]_g^+ z^\alpha(x) = z^\alpha(x)[K]_g^+ \) but \( [K]_g^- z^\alpha(x) = z^{-\alpha}(x)[K]_g^- \). So, another form of the Fourier expansion is \( a(x) = \sum_{n \in Z_i} ([a_n]_g^+ + [a_n]_g^-) z^n(x) \). Then, we let
\[ \overline{S}_{Z_i}^g(a) = \sum_{n \in Z_i} ([a_n]_g^+ - [a_n]_g^-) e_n. \]

For \( K \in \mathfrak{A} \) it yields \( S_{Z_i}^g(aK) = S_{Z_i}^g(a)K \) and \( S_{Z_i}^g(Ka) = ([K]_g^+ 1_{Z_i} + [K]_g^- T_{Z_i}) S_{Z_i}^g(a) \), ie. it is right \( \mathfrak{A} \)-linear and left \( \mathfrak{A} \)-quasi-linear. Similarly, \( \overline{S}_{Z_i}^g(Ka) = K \overline{S}_{Z_i}^g(a) \) and \( \overline{S}_{Z_i}^g(aK) = \overline{S}_{Z_i}^g(a) ([K]_g^+ 1_{Z_i} + [K]_g^- T_{Z_i}) \), ie. it is left \( \mathfrak{A} \)-linear and right \( \mathfrak{A} \)-quasi-linear.

When appropriate, right \( \mathfrak{A} \)-invariant linear maps from \( \mathfrak{B}^j \) to \( \mathfrak{B}^i \) can be represented by \( Z_i \times Z_j \)-matrices acting of the left of the column vectors considered above. If the action is \( A \) then for the representing matrix \( A^g \) we have \( S_{Z_j}^g(Aa) = A^g \overline{S}_{Z_i}^g(a) \). Similarly, through the representation \( S_{Z_i}^g(a) \) such matrices can also be used to represent left \( \mathfrak{A} \)-linear actions form \( \mathfrak{B}^j \) to \( \mathfrak{B}^i \).

b.) The problem with the skew-complex Fourier bases above is that sometimes there are no skew-involutions \( g \) in \( \mathfrak{A} \) at all. For arithmetic purposes it is always possible to adjoin an imaginary element \( i \) to \( \mathfrak{A} \), but this might change the invariants associated to \( \mathfrak{A} \). Hence it is much more natural to use real Fourier series. In the real representations we use bases corresponding to \( e_n \leftrightarrow \sqrt{2} \cos nx \ (n > 0) \), \( e_0 \leftrightarrow 1 \), \( e_{-n} \leftrightarrow \sqrt{2} \sin nx \ (n > 0) \). Our formulas using expansions in terms of \( z(x) = \cos x + g \sin x \) can be transposed to the real base with the help of the matrices
\[ N_{Z_i}^{l,g} = e_{0,0} + \frac{1}{\sqrt{2}} \sum_{n \in Z_i^+} -ge_{-n,-n} + ge_{-n,n} + e_{n,-n} + e_{n,n}, \]
\[ N_{Z_i}^{r,g} = \frac{1}{\sqrt{2}} \sum_{n \in Z_i^+} -ge_{-n,-n} + ge_{-n,n} + e_{n,-n} + e_{n,n}, \]
and \( N_{Z_i}^{g} = (N_{Z_i}^{l,g})^{-1} = (N_{Z_i}^{r,g})^\top \). Indeed, using the superscript “\( \top \)” in order to indicate real expansion we find that \( S_{Z_i}^l(a) = N_{Z_i}^{l,g} \overline{S}_{Z_i}^g(a) \). Similarly for \( S_{Z_i}^r(a) \). These real representations are plainly left and right \( \mathfrak{A} \)-linear. Furthermore, if \( A^g \) was a \( Z_i \times Z_j \) matrix representing an action then the corresponding matrix in the real basis yields
\[ A^r = N_{Z_i}^{l,g} A^g (N_{Z_j}^{r,g})^{-1}. \]

In practice, it is easier to carry out computations in the skew-complex representation, especially with \( g = i \), then translate the result to the real case. If the computation was not clear for \( g \neq i \) then we can still obtain the corresponding formula via the real case.

1.4. Multiplication actions. Suppose that the smooth function \( x \mapsto a(x) \) is in \( \mathfrak{B}^j \) and it has Fourier expansion \( a(x) = \sum_{n \in Z_i} z^n(x)a_n \). Suppose that \( Z_i = Z_l + Z_j \), then
\( a\mathcal{B}^j \subset \mathcal{B}^i \). Now, as \( a(x)z^k(x) = \sum_{n \in \mathbb{Z}} z^{n+k}(x)[a_n]^+ + z^{n-k}(x)[a_n]^g \) we see that multiplication by \( a(x) \) on the left can be represented by the matrix

\[
U^g_{z_i, z_j}(a) = \sum_{n \in \mathbb{Z}, k \in \mathbb{Z}} [a_n]^+ e_{n+k} + [a_n]^g e_{n-k} = \sum_{n \in \mathbb{Z}} [a_n]^g Z_{z_i, z_j}^g + [a_n]^g Z_{z_i, z_j}^g \mathbf{T}_{z_i, z_j},
\]

We see that \( S^g_{z_i, a}(a) = U^g_{z_i, z_j}(a) e_0 \), \( S^g_{z_j, a}(a) = e_0^T U^g_{z_j, z_i}(a) \). These skew-complex multiplication representations are \( \mathfrak{A} \)-quasi-linear in the sense that for \( K \in \mathfrak{A} \)

\[
U^g_{z_i, z_j}(K a) = ([K]^g_1 z_i + [K]^g_1^{-1} z_j) U^g_{z_i, z_j}(a),
\]

\[
U^g_{z_i, z_j}(a K) = U^g_{z_i, z_j}(a) ([K]^g_1 z_j + [K]^g_1^{-1} z_j) T_{z_i, z_j}
\]

Furthermore, the representations above are multiplicative: i.e. the identity

\[
U^g_{z_i, z_j}(a_1) U^g_{z_j, z_k}(a_2) = U^g_{z_i, z_k}(a_1 a_2)
\]

holds. Another simple but very useful observation which we call “real multiplication commutativity” is that for a periodic real scalar function \( b \) we know that

\[
U^g_{z_i, z_j}(a) U^g_{z_j, z_i}(b) = U^g_{z_j, z_i}(b) U^g_{z_i, z_j}(a).
\]

Everything can be transcribed to the real base, yielding various matrices \( U^g_{z_i, z_j}(a) \). These representations are plainly \( \mathfrak{A} \)-linear.

1.5. **Convention.** If we use indices \( i, j, l \) then it will be understood that \( \mathbb{Z}_i = \mathbb{Z}_i + \mathbb{Z}_j \).

1.6. The orientation reversal action \( C \) is defined by \( (x \mapsto a(x)) \mapsto (x \mapsto a(\bar{x})) \).

If \( a \in \mathcal{B}^i \) and \( a(x) = \sum_{n \in \mathbb{Z}} z^n (x) a_n \) then \( C a(x) = \sum_{n \in \mathbb{Z}} z^{-n}(x) a_n \). Its matrix is \( C^g_{z_i} = \mathbf{T}_{z_i} \) in the skew-complex base, and \( C^g_{z_i} = -1 \mathbf{1}_{z_i} + \mathbf{1}_{z_i} ^{-1} \) in the real base.

1.7. **Convention.** Instead of a pair of indices “\( g, x_1 \)” we may write “\( z_1 \)” where \( z_1 = \cos x_1 + g \sin x_1 \). We do this only when there is no danger of confusion.

1.8. **\( \delta \)-functions.** Let \( x_1 \) be a constant or a variable independent from \( x \). They are not smooth functions but only distributions, nevertheless we can consider the periodic and skew-periodic “algebraic” delta functions

\[
\delta^{x_1}_z(x) = 2\pi \sum_{n \in \mathbb{Z}} \delta_{x_1 + 2\pi n}(x) \quad \text{and} \quad \delta^{x_1}_z(x) = 2\pi \sum_{n \in \mathbb{Z}} (-1)^n \delta_{x_1 + 2\pi n}(x).
\]

The column and row matrices representing these distributions are

\[
\delta^{g,x_1}_{z_i} = \delta^{z_i}_z = \sum_{n \in \mathbb{Z}} z_1^{-n} e_n, \quad \delta^{g,x_1}_{z_i} = \delta^{z_i}_z = \sum_{n \in \mathbb{Z}} z_1^{-n} e_n^T.
\]

The real-based matrices are

\[
\delta^{x_1}_z = \sum_{n \in \mathbb{Z}^+} e_n \sqrt{2} \sin nx_1 + e_0 + \sum_{n \in \mathbb{Z}^+} e_n \sqrt{2} \cos nx_1, \quad \delta^{x_1}_z = (\delta^{x_1}_z)^T,
\]

\[
\delta^{x_1}_z = \sum_{n \in \mathbb{Z}^+} e_n \sqrt{2} \sin nx_1 + \sum_{n \in \mathbb{Z}^+} e_n \sqrt{2} \cos nx_1, \quad \delta^{x_1}_z = (\delta^{x_1}_z)^T.
\]

1.9. **Rotation actions.** Suppose that \( x_1 \) is a constant or a variable independent from \( x \). The rotation action \( R^{x_1} \) will be given as \( (x \mapsto a(x)) \mapsto (x \mapsto a(x - x_1)) \) for us.

We set \( z_1 = \cos x_1 + g \sin x_1 \). We define the rotation matrices as

\[
R^{g,x_1}_{z_i} = R^{z_i}_{x_1} = \sum_{n \in \mathbb{Z}} z_1^{-n} e_n, n.
\]
Indeed, $R_{z_i}^{g,x_1}S_{z_i}^g(a) = S_{z_i}^g(R^{x_1}a)$, and $S_{z_i}^g(a)(R_{z_i}^{g,x_1})^{-1} = S_{z_i}^g(R^{x_1}a)$. More generally,
\[ R_{z_i}^{g,x_1}U_{z_i,x_j}^g(a)(R_{z_j}^{g,x_1})^{-1} = U_{z_i,x_j}^g(R^{x_1}a). \]

Also, $R_{z_i}^{g,x_1}\delta_{z_i}^{g,x_2} = \delta_{z_i}^{g,x_1+x_2}$, $\delta_{z_i}^{g,x_1}(R_{z_i}^{g,x_2})^{-1} = \delta_{z_i}^{g,x_1+x_2}$. The real-based matrices are
\[ R_{z_i}^{r,x_1} = e_{0,0} + \sum_{n \in \mathbb{Z}^+} e_{-n,-n} \cos nx_1 + e_{-n,n} \sin nx_1, \]
\[ R_{z_i}^{r,x_1} = \sum_{n \in \mathbb{Z}^+} e_{-n,-n} \cos nx_1 + e_{-n,n} \sin nx_1. \]

1.10. Principal values of simple poles. Also quite important distributions are the distributions of Cauchy principal values of simple poles. The periodic versions are
\[ h_{z_i}^{x_1}(x) = \text{PV} \cot \frac{x-x_1}{2} \quad \text{and} \quad h_{z_i}^{x_1}(x) = \text{PV} \cosec \frac{x-x_1}{2}. \]

The column and row matrices representing these distributions are
\[ h_{z_i}^{g,x_1} \equiv h_{z_i}^{x_1} = \sum_{n \in \mathbb{Z}_i} -g(\text{sgn } n)z_i^{-n}e_n, \quad h_{z_i}^{g,x_1} \equiv h_{z_i}^{x_1} = \sum_{n \in \mathbb{Z}_i} -g(\text{sgn } n)z_i^n e_n^\top. \]

The real-based matrices are:
\[ h_{z_i}^{r,x_1} = \sum_{n \in \mathbb{Z}_i^+} e_{-n} \sqrt{2} \cos nx_1 - \sum_{n \in \mathbb{Z}_i^+} e_n \sqrt{2} \sin nx_1, \quad h_{z_i}^{r,x_1} = -(h_{z_i}^{r,x_1})^\top. \]

1.11. Integration and Hilbert transforms. For $u \in \mathfrak{A}$ the matrices
\[ f_{z_i}^g(u) = f_{z_i}^r(u) = u e_{0,0} \]
represent the “linear integration” operation $f_z : a \mapsto u \int_{y=0}^{2\pi} a(y) \frac{dy}{2\pi}$ on $\mathfrak{B}^0$. Let
\[ H_{z_i}^g = \sum_{n \in \mathbb{Z}_i} -g(\text{sgn } n)e_{n,n} = g1_{z_i} - g1_{z_i}^+, \]
\[ H_{z_i}^g[G] = g1_{z_i}^- + Ge_{0,0} - g1_{z_i}^+. \]

Then $H_{z_i}^g$ and $H_{z_i}^r$ are the matrices of the Hilbert transforms $H_{z_i}$ and $H_{z_i}$ on $\mathfrak{B}^0$ and $\mathfrak{B}^1$, respectively. Note that $H_{z_i}$ is not invertible while $H_{z_i}$ is invertible. If $G$ is a skew-involution then a reasonable invertible substitute for the even Hilbert transform is $H_{z_i}^g[G]$. The real-based matrices are
\[ H_{z_i}^g = \sum_{n \in \mathbb{Z}_i} \text{sgn } n e_{-n,n} \quad \text{and} \quad H_{z_i}^g[G] = Ge_{0,0} + \sum_{n \in \mathbb{Z}} \text{sgn } n e_{-n,n}. \]

1.12. Trivial extensions of loops. Let $z_1 = \cos x_1 + g \sin x_1$ again, and $u$ arbitrary. We define the matrices
\[ L_{z_i}^{g,x_1} = z_1 1_{z_i^-} + e_{0,0} + z_1^{-1}1_{z_i^+} = e_{0,0} + 1_{z_i\setminus\{0\}} \cos x_1 + H_{z_i}^g \sin x_1, \]
\[ L_{z_i}^{r,x_1} = z_1 1_{z_i^-} + z_1^{-1}1_{z_i^+} = 1_{z_i} \cos x_1 + H_{z_i}^r \sin x_1, \]
\[ L_{z_i}^{g,x_1} [u] = z_1 1_{z_i^-} + u e_{0,0} + z_1^{-1}1_{z_i^+} = u e_{0,0} + 1_{z_i\setminus\{0\}} \cos x_1 + H_{z_i}^g \sin x_1. \]

If $u$ is a general loop with variable $x_1$ then $L_{z_i}^{g,x_1} [u] \equiv L_{z_i}^{r,x_1} [u]$ can be considered as the trivial extension of $u$ by infinitely many space dimensions. The real-based matrices can quickly be recovered from the various constituents.
1.13. Up to this point we have defined several matrices $M^g$. We have seen that these matrices allow real-based transcriptions $M^g$ which do not depend on any skew-involution $g$, and, in fact, they do not require the existence of any such skew-involution. Colloquially, $M^g$ allows a real version $M^g$. In what follows we use the superscript “$g$” in order to indicate that such a transcription is possible, even if we do not elaborate it in detail. Then, it goes without saying that any formula involving various matrices $M^g$, linear combinations and multiplications allows a similar formula involving corresponding matrices $M^g$.

1.14. Algebraic kernel functions. For a rapidly decreasing $\mathbb{Z}_i \times \mathbb{Z}_j$ matrix

$$A^g = \sum_{n \in \mathbb{Z}_i, m \in \mathbb{Z}_j} a_{n,m} e_{n,m}$$

we define the its algebraic kernel function as

$$A(x_1, x_2) = \delta_{z_1}^{x_1} A^g \delta_{z_2}^{x_2} = \sum_{n \in \mathbb{Z}_i, m \in \mathbb{Z}_j} z_1^n a_{n,m} z_2^{-m}.$$ 

If $A^F$ is the real-based transcription of $A^g$ then $A(x_1, x_2) = \overline{\delta_{z_1}^{x_1} A^F \delta_{z_2}^{x_2}}$. In particular, if $A(x_1, x_2)$ is an expression independent of $g$ then $A^g$ allows a natural real transcription. In fact, instead of using (and examining) real matrices it is simpler to use the algebraic kernel function $A(x_1, x_2)$. Transformation kernels can also be defined, more generally, for matrices of temperate growth but then they yield distributions.

1.15. Grassmannians. The loop space skew-Grassmannians are defined by

$$Gr^g_{\mathbb{Z}} \mathcal{B}^0 = \{ U^g_{\mathbb{Z}}(a) H^g_{\mathbb{A}}(a^{-1}) : a \in \mathcal{B}^0 \text{ is invertible} \},$$

$$Gr^g_{\mathbb{Z}} \mathcal{B}^1 = \{ U^g_{\mathbb{Z},\mathbb{A}}(a) H^g_{\mathbb{A},\mathbb{Z}}(a^{-1}) : a \in \mathcal{B}^1 \text{ is invertible} \}.$$

We use the notations $J^g_{\mathbb{A}}(a)$ and $J^g_{\mathbb{Z}}(a)$ for the skew-involutions associated to the various invertible loops $a$ as above, respectively.

1.16. Theorem. $J^g_{\mathbb{A}}$ can be “inverted” by the formula

$$a(x_1) a(x_2)^{-1} = \overline{\delta_{z_1}^{x_1} \left( U^g_{\mathbb{Z},\mathbb{A}}(\sin \frac{x}{2}) J^g_{\mathbb{A}}(a) U^g_{\mathbb{Z},\mathbb{A}}(\cos \frac{x}{2}) - U^g_{\mathbb{Z},\mathbb{A}}(\cos \frac{x}{2}) J^g_{\mathbb{A}}(a) U^g_{\mathbb{Z},\mathbb{A}}(\sin \frac{x}{2}) \right) \delta_{z_2}^{x_2}}.$$

Similar formula holds with $J^g_{\mathbb{Z}}$ but with $\mathbb{Z}$ and $\mathbb{A}$ interchanged.

Proof. For $g = i$ the inversion formula

$$a(x_1) a(x_2)^{-1} = \left( \sum_{n \in \mathbb{Z}} z_1^n e_n^+ \right) U^g_{\mathbb{Z}}(a) e_{0,0} U^g_{\mathbb{Z}}(a^{-1}) \left( \sum_{m \in \mathbb{Z}} z_2^{-m} e_m \right) =$$

$$= \overline{\delta_{z_1}^{x_1} \left( U^g_{\mathbb{Z},\mathbb{A}}(z^{1/2}) J^g_{\mathbb{A}}(a) U^g_{\mathbb{Z},\mathbb{A}}(z^{-1/2}) - U^g_{\mathbb{Z},\mathbb{A}}(z^{-1/2}) J^g_{\mathbb{A}}(a) U^g_{\mathbb{Z},\mathbb{A}}(z^{1/2}) \right) \delta_{z_2}^{x_2}}$$

is easy to check. Expanded further, it yields our result, which form remains valid even after changing bases.

1.17. Theorem. The following are equivalent:

i.) $J \in Gr^g_{\mathbb{Z}} \mathcal{B}^1$

ii.) $J$ is a rapidly decreasing perturbation of $H^g_{\mathbb{Z}}$, $J^2 = -1_{\mathbb{Z}}$, and

$$U^g_{\mathbb{Z}}(\cos x) + JU^g_{\mathbb{Z}}(\cos x) J = U^g_{\mathbb{Z}}(\sin x) J - JU^g_{\mathbb{Z}}(\sin x).$$
Similar statement holds for $J \in G_{\mathfrak{A}}^2 \mathcal{B}^0$.

**Proof.** Assume that $g = i$. Now, it is not hard to see that (i) holds if and only if

$$\frac{1-iJ}{2} U^g_z(\frac{1+iJ}{2}) f g = \frac{1+iJ}{2} U^g_z(\frac{1-iJ}{2}) f g = 0$$

hold. (In an unprecise but possibly familiar language this is just the requirement that the $(-i)$-eigenspace of $J$ is $z$-invariant for multiplication, and the $i$-eigenspace of $J$ is $z^{-1}$-invariant for multiplication.) Expanding the two conditions in terms of $z = \cos x + i \sin x$ we obtain, somewhat surprisingly, a single equation as in (ii) for the case $g = i$. Changing bases the result follows in full generality. □

1.18. **Lemma.** If $G$ is a skew-involution then

$$H^g_{\bar{Z}}[G] = J^g_{\bar{Z}}(\cos \frac{x}{2} + G \sin \frac{x}{2}).$$

1.19. **Poisson and Hilbert-Poisson transforms.** Suppose that $a = \sum_{n \in \mathbb{Z}} z^n a_n$. For $-1 \leq r \leq 1$ we define the Poisson and Hilbert-Poisson transforms by

$$P^r_{\mathbb{A}} a = \sum_{n \in \mathbb{A}} r^{\lceil n \rceil} z^n a_n, \quad H^r_{\mathbb{A}} a = \sum_{n \in \mathbb{A}} - (\text{sgn } n) r^{\lceil n \rceil} z^n a_n.$$  

If $a = \sum_{n \in \mathbb{A}} z^n a_n$ then for $0 \leq r \leq 1$ we define

$$P^r_{\mathbb{A}} a = \sum_{n \in \mathbb{A}} r^{\lceil n \rceil} z^n a_n, \quad H^r_{\mathbb{A}} a = \sum_{n \in \mathbb{A}} - (\text{sgn } n) r^{\lceil n \rceil} z^n a_n.$$  

In the case $r = 1$ these operations specialize as follows: $P^1_{\mathbb{A}}$ yields the identity, and $H^1_{\mathbb{A}}$ yields the Hilbert transform.

2. **Some decompositions of matrices and related material**

2.1. In case of $L^g_{\bar{Z},Z}[u]$ the matrix was composed of two different parts. We can generalize this construction as follows. We define the matrices

$$\widetilde{H}^g_{\mathbb{A},\mathbb{Z}} = \sum_{n \in \mathbb{A}^-} g e_{\mathbb{A},n} - \sum_{n \in \mathbb{A}^+} g e_{\mathbb{A},n}, \quad \widetilde{H}^g_{\mathbb{Z},\mathbb{A}} = (\widetilde{H}^g_{\mathbb{A},\mathbb{Z}})^\top,$$

$$\widetilde{H}^g_{\mathbb{Z},\mathbb{A}} = \sum_{n \in \mathbb{Z}^+} g e_{\mathbb{Z},-\frac{1}{2},n} + \frac{\sqrt{2}}{2} e_{\mathbb{Z},\frac{1}{2},0} - e_{\mathbb{Z},0} - \sum_{n \in \mathbb{Z}^-} g e_{\mathbb{Z},\frac{1}{2},n}, \quad \widetilde{H}^g_{\mathbb{A},\mathbb{Z}} = (\widetilde{H}^g_{\mathbb{Z},\mathbb{A}})^\top;$$

and we let $f^g_{\mathbb{A},i} = e_{0,i}, f^g_{\mathbb{A},i} = \sqrt{2} e_{\frac{1}{2},i} + \sqrt{2} e_{\frac{1}{2},i}, f^g_{\mathbb{Z},i} = e_{0,i}, f^g_{\mathbb{Z},i} = \sqrt{2} e_{\frac{1}{2},i} + \sqrt{2} e_{\frac{1}{2},i}$.

If $A$ is a $\mathbb{Z}_i \times \mathbb{Z}_j$ matrix, $B$ is a $\mathbb{Z}_i$-column matrix, $C$ is a $\mathbb{Z}_j$-row matrix, $d$ is a scalar matrix, then we define

$$\begin{bmatrix} A & B \\ C & d \end{bmatrix}^g_{\mathbb{Z}_i,\mathbb{Z}_j} = \widetilde{H}^g_{\mathbb{Z}_i,\mathbb{Z}_j} A \widetilde{H}^g_{\mathbb{Z}_j,\mathbb{Z}_i} + \widetilde{H}^g_{\mathbb{Z}_i,\mathbb{Z}_i} B \otimes f^g_{\mathbb{Z}_j} + f^g_{\mathbb{Z}_i} \otimes C \widetilde{H}^g_{\mathbb{Z}_j,\mathbb{Z}_i} + df^g_{\mathbb{Z}_i} \otimes f^g_{\mathbb{Z}_j}.$$  

These symbols multiply as $2 \times 2$ matrices. Every $\mathbb{Z}_i \times \mathbb{Z}_j$ matrix $C$ occurs so:

$$C = \begin{bmatrix} \widetilde{H}^g_{\mathbb{Z}_i,\mathbb{Z}_i} C \widetilde{H}^g_{\mathbb{Z}_j,\mathbb{Z}_j} & \widetilde{H}^g_{\mathbb{Z}_i,\mathbb{Z}_j} C f^g_{\mathbb{Z}_j} \\ f^g_{\mathbb{Z}_i} C \widetilde{H}^g_{\mathbb{Z}_j,\mathbb{Z}_j} & f^g_{\mathbb{Z}_i} C f^g_{\mathbb{Z}_j} \end{bmatrix}^g_{\mathbb{Z}_i,\mathbb{Z}_j}.$$
There are variants of this construction but without \( \frac{\sqrt{2}}{2} \). Let \( X^g_Z = 1_Z \), and
\[
X^g_Z = 1_{Z \setminus \{-\frac{1}{2}, \frac{1}{2}\}} + \frac{\sqrt{2}(e_{-\frac{1}{2}} - e_{\frac{1}{2}}) \otimes (e^T_{-\frac{1}{2}} - e^T_{\frac{1}{2}})}{2} + \frac{\sqrt{2}(e_{-\frac{1}{2}} + e_{\frac{1}{2}}) \otimes (e^T_{-\frac{1}{2}} + e^T_{\frac{1}{2}})}{2};
\]
and we define
\[
\widetilde{H}^{R,g}_{Z_i, Z_i} = \bar{X}^g_{Z_i} \widetilde{H}^{g}_{Z_i, Z_i}, \quad \mathbf{f}^{R,g}_{Z_i} = X^g_{Z_i} \mathbf{f}^g_{Z_i}, \quad \widetilde{H}^{R,g}_{Z_i, Z_i} = \bar{X}^g_{Z_i} \mathbf{f}^{g}_{Z_i} - \mathbf{f}^{g}_{Z_i}, \quad \mathbf{f}^{R,g}_{Z_i} = \bar{X}^g_{Z_i} \mathbf{f}^{g}_{Z_i} - \mathbf{f}^{g}_{Z_i} \mathbf{f}^{g}_{Z_i}.
\]
Practically this means that every term \( \frac{\sqrt{2}}{2} \) is replaced by 1 or \( \frac{1}{2} \).

Then we can define \( \{A, b\}_{R,g} \) and \( \{A, b\}_{L,g} \) analogously. In fact, the relationship between the basic version and these is very simple:
\[
\{A, b\}_{R,g} = \mathbf{X}^g_{Z_i} \{A, b\}^{g}_{Z_i, Z_i}, \quad \{A, b\}_{L,g} = \mathbf{X}^g_{Z_i} \{A, b\}^{g}_{Z_i, Z_i}.
\]
The relationship to our original motivating example is that
\[
L^{g,x;1}_{Z}[u] = \left\{ \begin{array}{c} L^{g,x;1}_{Z} u \end{array} \right\}_{Z} = \left\{ \begin{array}{c} \frac{g, x; 1}{\bar{Z}} u \end{array} \right\}_{Z} = \left\{ \begin{array}{c} \frac{g, x; 1}{Z} u \end{array} \right\}_{Z}.
\]
In general, \( \{A, b\}_{Z} \) and \( \{d_1 1_{Z_i}, d_2 \}_{Z_i} \) do not depend on the choice of \( X \), which can be \( \emptyset \) (the empty symbol), \( R \) or \( L \).

2.2. Difference functions. For \( a \in \mathcal{B}^0 \) we consider the difference functions
\[
\Delta^{x_1}_{Z,Z} a(x) = \frac{1}{2} \left( a(x) - a(x_1) \right) \cos \frac{x - x_1}{2},
\]
\[
\Delta^{x_1}_{Z,Z} a(x) = \frac{1}{2} \left( a(x) - a(x_1) \right) \sin \frac{x - x_1}{2},
\]
and for \( a \in \mathcal{B}^1 \)
\[
\Delta^{x_1}_{Z,Z} a(x) = \frac{1}{2} \left( a(x) - a(x_1) \right) \cos \frac{x - x_1}{2},
\]
\[
\Delta^{x_1}_{Z,Z} a(x) = \frac{1}{2} \cos \frac{x - x_1}{2},
\]
2.3. Lemma. For \( n \in \mathbb{Z}^+ \) it yields
\[
\Delta^{x_1}_{Z,Z} (z^n) = -g \left( \frac{1}{2} z^n + z_1^{-1} z^{n-1} + \ldots + z_1^{-n+1} z^{-1} + \frac{1}{2} z_1^{n-1} \right),
\]
\[
\Delta^{x_1}_{Z,Z} (1) = 0
\]
\[
\Delta^{x_1}_{Z,Z} (z^n) = g \left( \frac{1}{2} z^n + z_1^{-1} z + \ldots + z_1 z^{n-1} + \frac{1}{2} z^n \right),
\]
Lemma. It yields

\[
\int \Delta_{\mathbb{Z},\mathbb{Z}}^1 a(x) \frac{dx}{2\pi} = -\frac{1}{2} H_{\mathbb{Z}} a(x_1), \quad \int \Delta_{\mathbb{Z},\mathbb{Z}}^2 a(x) \frac{dx}{2\pi} = -\frac{1}{2} H_{\mathbb{Z}} a(x).
\]

2.5. We define the critical part of the difference functions as

\[
\tilde{\Delta}_{\mathbb{Z},\mathbb{Z}}^1 a(x) = -\frac{1}{2} \delta_{1,0} H_{\mathbb{Z}} a(x_1) - \frac{1}{2} \delta_{j,0} H_{\mathbb{Z}} a(x).
\]

2.6. Multiplicative difference functions. If \(a\) is multiplicatively invertible then it is also natural to consider multiplicative difference functions: For \(a \in \mathfrak{B}^0\)

\[
\Xi_{\mathbb{Z},\mathbb{Z}} a(x, x_1) = \frac{1}{2} \left( a(x) a(x_1)^{-1} - 1 \right) \cos \frac{x - x_1}{2},
\]

\[
\Xi_{\mathbb{Z},\mathbb{Z}} a(x, x_1) = \frac{1}{2} \left( a(x) a(x_1)^{-1} - 1 \right) \sin \frac{x - x_1}{2},
\]

and for \(a \in \mathfrak{B}^1\)

\[
\Xi_{\mathbb{Z},\mathbb{Z}} a(x, x_1) = \frac{1}{2} \left( a(x) a(x_1)^{-1} - \cos \frac{x - x_1}{2} \right),
\]

\[
\Xi_{\mathbb{Z},\mathbb{Z}} a(x, x_1) = \frac{1}{2} \left( a(x) a(x_1)^{-1} - \sin \frac{x - x_1}{2} \right).
\]

2.7. Lemma.

\[
\int \Xi_{\mathbb{Z},\mathbb{Z}} a(x, x_1) \frac{dx}{2\pi} = -\frac{1}{2} H_{\mathbb{Z}} a(x_1) \cdot a(x_1)^{-1},
\]

\[
\int \Xi_{\mathbb{Z},\mathbb{Z}} a(x_1, x) \frac{dx}{2\pi} = \frac{1}{2} a(x_1) \cdot H_{\mathbb{Z}} (a^{-1})(x_1).
\]
2.8. We can define normalized multiplicative difference functions by

\[ \Xi_{z_i}^R a(x, x_1) = \Xi_{z_i}^L a(x, x_1) + \frac{1}{2} \delta_{i,0} H_{z_i} a(x_1) a(x_1)^{-1} + \frac{1}{2} \delta_{j,0} H_{z_i} a(x) a(x_1)^{-1}, \]

\[ \Xi_{z_i}^L a(x_1) = \Xi_{z_i} a(x_1) - \frac{1}{2} \delta_{i,0} H_{z_i} a^{-1}(x_1) - \frac{1}{2} \delta_{j,0} a(x_1) H_{z_i} a^{-1}(x). \]

Then

\[ \int \Xi_{z_i}^L a(x, x_1) \frac{dx}{2\pi} = 0, \quad \text{and} \quad \int \Xi_{z_i}^R a(x_1, x) \frac{dx}{2\pi} = 0. \]

2.9. The following table might help to keep track the various periodicity properties:

| \( l \) | \( i \) | \( j \) | \( a(x) \) in x | \( \Delta_{z_i} \Xi_{z_i}^L a(x) \) in x | \( \Xi_{z_i} a(x_1) \) in x or x_1 |
|-------|-------|-------|---------------|-----------------|------------------|
| 0     | 0     | 0     | per           | \( \Xi_{z_i} a(x_1) \) in x or x_1 | \( \Xi_{z_i} a(x_1) \) in x or x_1 |
| 0     | 1     | 1     | per           | skew            | skew             |
| 1     | 0     | 1     | skew          | per             | skew             |
| 1     | 1     | 0     | skew          | per             | per              |

2.10. Lemma. It yields

\[ \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
S_{z_i}^g (\Delta_{z_i} \Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} = \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \]

\[ = \left\{ \begin{array}{ll}
1_{z_i}^g \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
a(x_1)
\end{array} \right\}_{z_i, z_j} \]

\[ = \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \left\{ \begin{array}{ll}
1_{z_i}^g \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \]

\[ = \left\{ \begin{array}{ll}
1_{z_i}^g \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
a(x_1)
\end{array} \right\}_{z_i, z_j} \]

\[ = \left\{ \begin{array}{ll}
1_{z_i}^g \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
a(x_1)
\end{array} \right\}_{z_i, z_j} \]

and

\[ = \left\{ \begin{array}{ll}
1_{z_i}^g \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
a(x_1)
\end{array} \right\}_{z_i, z_j} \]

\[ = \left\{ \begin{array}{ll}
1_{z_i}^g \\
S_{z_i}^g (\Xi_{z_i}^L a(x_1))
\end{array} \right\}_{z_i, z_j} \left\{ \begin{array}{ll}
U_{z_i, z_j}^g (a) \\
a(x_1)
\end{array} \right\}_{z_i, z_j} \]
3.1. For $n \in \frac{1}{2} \mathbb{Z}$ and a “scalar” $A$ we let $I_n(A)$ be the matrix which is $1_{(Z+n)}$ but the coefficient of $e_{n,n}$ is replaced by $A$, i.e.

$$I_n(A) = 1_{(Z+n)} \setminus \{n\} + Ae_{n,n}.$$ 

For $n \in \frac{1}{2} \mathbb{Z}$ and a $2 \times 2$ matrix $A$ we let $\mathcal{I}_n(A)$ be the matrix which is $1_{(Z+n)}$ but the $\{n - \frac{1}{2}, n + \frac{1}{2}\} \times \{n - \frac{1}{2}, n + \frac{1}{2}\}$ block is replaced by $A$. More precisely, we set

$$\mathcal{I}_n\left(\begin{bmatrix} a & b \\ c & d \end{bmatrix}\right) = 1_{(Z+n) \setminus (\{n - \frac{1}{2}, n + \frac{1}{2}\} \times \{n - \frac{1}{2}, n + \frac{1}{2}\})} + a e_{n - \frac{1}{2}, n - \frac{1}{2}} + b e_{n - \frac{1}{2}, n + \frac{1}{2}} + c e_{n + \frac{1}{2}, n - \frac{1}{2}} + d e_{n + \frac{1}{2}, n + \frac{1}{2}}.$$ 

3.2. We define

$$A_n(\theta) = \mathcal{I}_n\left(\begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}\right), \quad A_n(\theta) = \mathcal{I}_n\left(\begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}\right),$$

where $\theta \in \mathbb{R}$. Similarly, we can define

$$A_n^R(\theta) = \mathcal{I}_n\left(\begin{bmatrix} 1 & \sin \theta \\ -\sin \theta & 1 - \sin^2 \theta \end{bmatrix}\right), \quad A_n^R(\theta) = \mathcal{I}_n\left(\begin{bmatrix} 1 & \sin \theta \\ -\sin \theta & 1 \end{bmatrix}\right),$$

$$A_n^L(\theta) = \mathcal{I}_n\left(\begin{bmatrix} 1 - \sin^2 \theta & \sin \theta \\ -\sin \theta & 1 \end{bmatrix}\right), \quad A_n^L(\theta) = \mathcal{I}_n\left(\begin{bmatrix} 1 & \sin \theta \\ -\sin \theta & 1 \end{bmatrix}\right).$$
In this case we can use natural parametrization by \( t = \sin \theta \). We present yet another version, which is less nice than the previous ones but it may illustrate the structure of these matrices in question a bit further. We define

\[
A_{n}^{R} (\theta) = \eta(n) \left( \begin{bmatrix} \cos \theta \cos \frac{\theta}{3} & \sin \theta \\ -\sin \theta & \cos \frac{\theta}{3} \end{bmatrix} \right), \quad A_{n}^{T} (\theta) = \eta(n) \left( \begin{bmatrix} \cos \frac{2\theta}{3} & \sin \frac{2\theta}{3} \\ -2\sin \frac{\theta}{3} & 1 \end{bmatrix} \right),
\]

\[
A_{n}^{L} (\theta) = \eta(n) \left( \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \frac{\theta}{3} \end{bmatrix} \right), \quad A_{n}^{\bar{L}} (\theta) = \eta(n) \left( \begin{bmatrix} 1 & 2\sin \frac{\theta}{3} \\ -\sin \theta & \cos \frac{2\theta}{3} \end{bmatrix} \right).
\]

In this case we can use natural parametrization by \( \tau = 2\sin \frac{\theta}{3} \). Then

\[
\sin \frac{\theta}{3} = \frac{\tau}{2}, \quad \cos \frac{2\theta}{3} = 1 - \frac{\tau^2}{2}, \quad \cos \theta \cos \frac{\theta}{3} = 2\cos \frac{2\theta}{3} - 1 = 1 - \tau^2,
\]

\[
\cos \theta \cos \frac{\theta}{3} = \frac{1}{2} \left( \cos \frac{2\theta}{3} + \cos \frac{4\theta}{3} \right) = \frac{(1 - \tau^2)(4 - \tau^2)}{4}, \quad \sin \theta = \frac{\tau(3 - \tau^2)}{2};
\]

so the entries of the replacement blocks are simple expressions of \( \tau \).

3.3. We will use the symbol \( X \) for any of \( R, R, \emptyset \) (the empty symbol), \( \frac{R}{3}, L \). Then the symbol \( -X \) will be used for \( L, \frac{R}{3}, \emptyset, \frac{R}{3}, R \), respectively.

The matrices \( A_{n}^{X} (\theta) \) are to be considered as elementary rotations, and the matrices \( \bar{A}_{n}^{X} (\theta) \) are to be considered as elementary half-rotations. It yields

\[
A_{n}^{X} (\theta) = \bar{A}_{n}^{X} (\theta) A_{n}^{-X} (-\theta)^{-1}, \quad A_{n}^{X} (\theta)^\top = \bar{A}_{n}^{-X} (\theta)^{-1}, \quad T_{Z_{n}} A_{n}^{X} (\theta) T_{Z_{n}} = \bar{A}_{n}^{-X} (\theta);
\]

hence

\[
A_{n}^{X} (\theta) = A_{n}^{-X} (-\theta)^{-1}, \quad A_{n}^{X} (\theta)^\top = A_{n}^{X} (-\theta), \quad T_{Z_{n}} A_{n}^{X} (\theta) T_{Z_{n}} = A_{n}^{X} (-\theta)^{-1};
\]

although for particular choices of \( X \) more special equalities hold.

In the rest of the section we assume that \( \theta \notin \pi Z \), i.e. that \( \sin \theta \neq \pm 1 \).

3.4. In what follows, it is useful to consider the function

\[
\omega(x, \theta) \equiv \omega(x, z^{-1}, \theta) = (1 - z \sin \theta)(1 - z^{-1} \sin \theta) = 1 - 2\cos x \sin \theta + \sin^2 \theta;
\]

even if only as an abbreviation. It is a positive real function. In particular, we can take various powers of it without the problem of ambiguities.

3.5. Let us consider the matrices \( \bar{A}_{Z}^{R,R,g} (\theta) = 1_{Z} \),

\[
\bar{A}_{Z}^{R,R,g} (\theta) = l_0 \left( \begin{bmatrix} \cos \theta & \cos \frac{\theta}{3} \\ -\sin \theta & \cos \frac{\theta}{3} \end{bmatrix} \right), \quad \bar{A}_{Z}^{R,R,g} (\theta) = l_0(\cos \theta),
\]

\[
\bar{A}_{Z}^{L,R,g} (\theta) = l_0 \left( \begin{bmatrix} \cos \theta & \cos \frac{\theta}{3} \\ -\sin \theta & \cos \frac{\theta}{3} \end{bmatrix} \right), \quad \bar{A}_{Z}^{L,R,g} (\theta) = l_0(\cos \theta);
\]

and \( \bar{A}_{Z}^{R,R,g} (\theta) = 1_{Z} \),

\[
\bar{A}_{Z}^{R,R,g} (\theta) = l_0 \left( \begin{bmatrix} \cos \frac{2\theta}{3} & -\sin \frac{2\theta}{3} \\ -\sin \frac{\theta}{3} & \cos \frac{2\theta}{3} \end{bmatrix} \right), \quad \bar{A}_{Z}^{R,R,g} (\theta) = l_0 \left( \begin{bmatrix} \cos \frac{\theta}{3} & -\sin \frac{\theta}{3} \\ -\sin \frac{\theta}{3} & \cos \frac{2\theta}{3} \end{bmatrix} \right),
\]

\[
\bar{A}_{Z}^{L,R,g} (\theta) = l_0 \left( \begin{bmatrix} 1 & -\sin \theta \\ -2\sin \frac{\theta}{3} & 1 \end{bmatrix} \right), \quad \bar{A}_{Z}^{L,R,g} (\theta) = l_0 \left( \begin{bmatrix} 1 & -\sin \theta \\ -\sin \theta & -\sin \theta \end{bmatrix} \right).
\]

One can extend this notation so that \( \bar{A}_{Z_{n}}^{X,Y,g} (\theta) = \bar{A}_{Z_{n}}^{X,Z,g} (\theta) \bar{A}_{Z_{n}}^{Z,Y,g} (\theta) \) would hold.
3.6. Consider the matrices

\[
F_{Z_i}^{R,g}(\theta) = 1_{Z_i} - U_{Z_i}^g (1 - z^{-1} \sin \theta) + e_{0,0} + 1_{Z_i} U_{Z_i}^g (1 - z \sin \theta) = \\
= 1_{Z_i} - U_{Z_i}^g (\cos x) \sin \theta + H_{Z_i}^g U_{Z_i}^g (\sin x) \sin \theta + e_{0,0} U_{Z_i}^g (\cos x) \sin \theta = \\
= 1_{Z_i} - U_{Z_i}^g (\cos x) \sin \theta + U_{Z_i}^g (\sin x) H_{Z_i}^g \sin \theta - e_{0,0} U_{Z_i}^g (\cos x) \sin \theta,
\]

\[
F_{X_i}^{R,g}(\theta) = 1_{X_i} - U_{X_i}^g (1 - z^{-1} \sin \theta) + 1_{(-\frac{1}{3}, \frac{1}{2})} + 1_{X_i} U_{X_i}^g (1 - z \sin \theta) = \\
= U_{X_i}^g (1 - z^{-1} \sin \theta) 1_{X_i} - U_{X_i}^g (1 - z \sin \theta) 1_{X_i} = \\
= 1_{Z_i} - U_{Z_i}^g (\cos x) \sin \theta + U_{X_i}^g (\sin x) H_{X_i}^g \sin \theta;
\]

and, furthermore, we define the matrices

\[
F_{Z_i}^{X,g}(\theta) = \overline{A}_{Z_i}^{X,R,g}(\theta) F_{Z_i}^{R,g}(\theta) = \overline{A}_{Z_i}^{X,R,g}(\theta) + F_{Z_i}^{R,g}(\theta) - 1_{Z_i}.
\]

In particular, it yields

\[
F_{Z_i}^{2,g}(\theta) = 1_{Z_i} - U_{Z_i}^g (1 - z^{-1} \sin \theta) + e_{0,0} \cos \theta + 1_{Z_i} U_{Z_i}^g (1 - z \sin \theta),
\]

\[
F_{X_i}^{2,g}(\theta) = 1_{X_i} - U_{X_i}^g (1 - z^{-1} \sin \theta) + e_{0,0} \cos \theta + 1_{X_i} U_{X_i}^g (1 - z \sin \theta),
\]

\[
F_{Z_i}^{L,g}(\theta) = 1_{Z_i} - U_{Z_i}^g (1 - z^{-1} \sin \theta) + e_{0,0} (1 - \sin^2 \theta) + 1_{Z_i} U_{Z_i}^g (1 - z \sin \theta),
\]

\[
F_{X_i}^{L,g}(\theta) = 1_{X_i} - U_{X_i}^g (1 - z^{-1} \sin \theta) + 1_{X_i} U_{X_i}^g (1 - z \sin \theta) = \\
= 1_{Z_i} - U_{Z_i}^g (\cos x) \sin \theta + H_{X_i}^g U_{X_i}^g (\sin x) \sin \theta.
\]

We also define

\[
F_{Z_i}^{X,g}(\theta) = F_{Z_i}^{X,g}(\theta)^{\top}.
\]

3.7. Lemma. It yields

\[
F_{Z_i}^{X,g}(\theta) U_{Z_i}^g (\omega(x, \theta)^{-1}) F_{Z_i}^{X,g}(\theta) = 1_{Z_i}.
\]

3.8. For \( c = -1, -\frac{1}{3}, 0, \frac{1}{3}, 1 \) let \( cL \) denote \( R, \frac{R}{3}, \emptyset, \frac{L}{3}, L \) respectively. Let

\[
D_{Z_i}^{cL,g}(\theta) = F_{Z_i}^{cL,g}(\theta) U_{Z_i}^g (\omega(x, \theta)^{-\frac{1}{cL}}).
\]

In particular, it yields

\[
D_{Z_i}^{R,g}(\theta) = 1_{Z_i} - U_{Z_i}^g (1 - z^{-1} \sin \theta) + e_{0,0} + 1_{Z_i} U_{Z_i}^g (1 - z \sin \theta),
\]

\[
D_{X_i}^{R,g}(\theta) = 1_{X_i} - U_{X_i}^g (1 - z^{-1} \sin \theta) + 1_{(-\frac{1}{3}, \frac{1}{2})} + 1_{X_i} U_{X_i}^g (1 - z \sin \theta),
\]

\[
D_{Z_i}^{R,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + e_{0,0} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_{Z_i}^{L,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + e_{0,0} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_{Z_i}^{L,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + 1_{Z_i} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_{Z_i}^{R,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + e_{0,0} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_{Z_i}^{L,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + 1_{Z_i} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_{Z_i}^{R,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + e_{0,0} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_{Z_i}^{L,g}(\theta) = 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z^{-1} \sin \theta}{1 - z \sin \theta} \right) + 1_{Z_i} U_{Z_i}^g \left( \frac{\cos \theta}{\sqrt{\omega(z, z^{-1}, \theta)}} \right) + \\
+ 1_{Z_i} U_{Z_i}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right),
\]
\[
D_z^L(\theta) = 1_{\mathbb{Z}} - U_{\mathbb{Z}}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right) + e_{0,0} U_{\mathbb{Z}}^g \left( \frac{1 - \sin^2 \theta}{\omega(z, z^{-1}, \theta)} \right) + 1_{\mathbb{Z}} + U_{\mathbb{Z}}^g \left( \frac{1}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_z^L(\theta) = 1_{\mathbb{Z}} - U_{\mathbb{Z}}^g \left( \frac{1}{1 - z \sin \theta} \right) + e_{0,0} U_{\mathbb{Z}}^g \left( \frac{1 - \sin^2 \theta}{\omega(z, z^{-1}, \theta)} \right) + 1_{\mathbb{Z}} + U_{\mathbb{Z}}^g \left( \frac{1}{1 - z^{-1} \sin \theta} \right),
\]

\[
D_z^L(\theta) = 1_{\mathbb{Z}} - U_{\mathbb{Z}}^g \left( \frac{1}{1 - z \sin \theta} \right) + e_{0,0} U_{\mathbb{Z}}^g \left( \frac{1}{1 - z^{-1} \sin \theta} \right).
\]

One can check that
\[
D_z^L(\theta)^{-1} = U_{\mathbb{Z}}^g(\omega(x, \theta)^{-\frac{1}{2}}) F_{\mathbb{Z}}^{cL}(\theta);
\]
or, in different terms,
\[
D_z^L(\theta)^{-1} = D_z^{-X,g}(\theta)^\top.
\]

Furthermore, \(T_{\mathbb{Z}}\) commutes with \(D_z^L(\theta)\).

3.9. The matrices \(D_z^L(\theta)\) and \(A_0^X(\theta)\) are closely related in the skew-complex representation. If we take sufficiently weak limits, entrywise, for example, then we find
\[
U_{\mathbb{Z}}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right) = \lim_{n,m \to \infty} A_{-m-\frac{1}{2}}^X(\theta) \ldots A_{-\frac{1}{2}}^X(\theta) A_{\frac{1}{2}}^X(\theta) \ldots A_{n+\frac{1}{2}}^X(\theta)
\]
and
\[
U_{\mathbb{Z}}^g \left( \frac{1 - z \sin \theta}{1 - z^{-1} \sin \theta} \right) = \lim_{n,m \to \infty} A_{-m}^X(\theta) \ldots A_0^X(\theta) \ldots A_n^X(\theta);
\]
where \(n, m \in \mathbb{Z}\). Similarly, it yields
\[
F_{\mathbb{Z}}^{X,g}(\theta) U_{\mathbb{Z}}^g \left( \frac{1}{1 - z^{-1} \sin \theta} \right) = \lim_{n \to \infty} A_{\frac{1}{2}}^X(\theta) \ldots A_{n+\frac{1}{2}}^X(\theta)
\]
and
\[
F_{\mathbb{Z}}^{X,g}(\theta) U_{\mathbb{Z}}^g \left( \frac{1}{1 - z^{-1} \sin \theta} \right) = \lim_{n \to \infty} A_{0}^X(-\theta)^{-1} A_{1}^X(\theta) \ldots A_n^X(\theta).
\]

Hence, by taking \(\sqrt{A \cdot B} = B \sqrt{B^{-1} A}\), we obtain
\[
F_{\mathbb{Z}}^{X,g}(\theta) U_{\mathbb{Z}}^g(\omega(x, \theta)^{-\frac{1}{2}}) = \sqrt{\lim_{m \to \infty} A_{-\frac{1}{2}}^X(\theta) \ldots A_{-m}^X(\theta)^{-1}) \cdot \lim_{n \to \infty} A_{\frac{1}{2}}^X(\theta) \ldots A_{n+\frac{1}{2}}^X(\theta)}
\]
and
\[
F_{\mathbb{Z}}^{X,g}(\theta) U_{\mathbb{Z}}^g(\omega(x, \theta)^{-\frac{1}{2}}) = \sqrt{\lim_{m \to \infty} A_{0}^X(-\theta)^{-1} A_{1}^X(\theta) \ldots A_n^X(\theta)}.
\]

Consequently, up to multiplication by \(U_{\mathbb{Z}}^g(\omega(x, \theta)^{-\frac{1}{2}})\) the matrices \(D_z^L(\theta)\) can be obtained from elementary rotations. The equalities above may also serve as
alternative definitions. In the cases $X = R, \frac{R}{2}, \frac{R}{3}, L$ one might even argue that $U_{\mathbb{Z}}^g(\omega(x, \theta)c)$ can be obtained from half-rotations.

3.10. Every matrix $M_{\mathbb{Z}i,\mathbb{Z}j}^g$ defined in this section can be twisted by rotations, so we obtain the matrices $M_{\mathbb{Z}i,\mathbb{Z}j}^g(x_1) = R_{\mathbb{Z}i}^g x_1 M_{\mathbb{Z}i,\mathbb{Z}j}^g R_{\mathbb{Z}j}^g(-x_1)$. For example, from $M_{\mathbb{Z}i}^g(\theta)$ one can prepare the matrices

$$M_{\mathbb{Z}i}^g(\theta, x_1) = R_{\mathbb{Z}i}^g x_1 M_{\mathbb{Z}i}^g(\theta) R_{\mathbb{Z}i}^g(-x_1).$$

In case of $U_{\mathbb{Z}i}^g(x \mapsto f(x))$ we obtain $U_{\mathbb{Z}i}^g(x \mapsto f(x-x_1))$ instead. Then every arithmetical relation which holds among these kind of matrices remains true for the extended matrices. (But not the ones which contain transposition, because it is not invariant.) So, in the computations above “$(\theta)$” should be replaced by “$(\theta, x_1)$”, and “$\omega(x, \theta)$” should be replaced by “$\omega(x-x_1, \theta)$”.

3.11. Lemma. It yields

$$F^X_{\mathbb{Z}i} X, g(\theta, x_1) \left( 1_{\mathbb{Z}i} \cos \frac{x - x_2}{2} + H_{\mathbb{Z}i}^g \sin \frac{x - x_2}{2} \right) F^X_{\mathbb{Z}i} \theta, x_2 =$$

$$= 1_{\mathbb{Z}i} \cos \frac{x - x_2}{2} (1 + \sin^2 \theta) + H_{\mathbb{Z}i}^g \sin \frac{x - x_2}{2} \cos^2 \theta - U_{\mathbb{Z}i}^g (2 \cos \frac{x - x_2}{2} - x) \sin \theta.$$

$\square$

4. Deformations of multiplication actions

4.1. Now we combine multiplication matrices and the special matrices introduced in the previous section. Let us define

$$U_{\mathbb{Z}i,\mathbb{Z}j}(a, \theta, x_1) = D_{\mathbb{Z}i}^g(\theta, x_1) U_{\mathbb{Z}i,\mathbb{Z}j}^g(a) D_{\mathbb{Z}j}^g(\theta, x_1)^{-1}.$$

Right now $U_{\mathbb{Z}i,\mathbb{Z}j}^g(a, \theta, x_2)$ is defined only for $\theta \notin \pi \mathbb{Z}$, i.e. for $\sin \theta \neq \pm 1$. But we will shortly see that the definition extends. Let us start with the case of $U_{\mathbb{Z}}^g(a, \theta, x_1)$.

One finds that

$$U_{\mathbb{Z}}^g \left( \sum_{n \in \mathbb{Z}} z^n a_n, \theta, x_1 \right) = \sum_{n \in \mathbb{Z}} U_{\mathbb{Z}}^g(z^n, \theta, x_1) \left( [a_n]^+ T_{\mathbb{Z}} + [a_n]^- T_{\mathbb{Z}} \right).$$

Using the abbreviations $s = \cos \theta$, $t = \sin \theta$, we see that for $n > 0$ it yields

$$U_{\mathbb{Z}}^g(z^n, \theta, x_1) =$$

$$\begin{bmatrix}
\vdots \\
1 \\
-\frac{z_1^{-1} t s z_1 s}{s^2} & \vdots & t z_1 s^{-1} s^{-1} s z_1^{-1} & \vdots \\
\vdots \\
-t s z_1 s^{-1} & \vdots & t z_1 s^{-1} & \vdots \\
\vdots \\
-t z_1^{-1} & \vdots & 1 & \vdots \\
\end{bmatrix}$$

and

$$U_{\mathbb{Z}}^g(1, \theta, x_1) = 1_{\mathbb{Z}};$$

furthermore

$$U_{\mathbb{Z}}^g(z^{-n}, \theta, x_1) = T_{\mathbb{Z}} U_{\mathbb{Z}}^g(z^n, \theta, -x_1) T_{\mathbb{Z}},$$
which means that $U^g_{z_1}(z^n, \theta, x_1)$ gets reflected and $z_1$ should be substituted by $z_1^{-1}$.

From the nature of these expressions we see that $U^g_{z_1}(a, \theta, x_1)$ extends to all $\theta$, and all along it yields a rapidly decreasing perturbation of $U^g_{z_1}(a) = U^g_{z_1}(a, 0, 0)$. The whole discussion extends to the general case. In fact, due to multiplicativity, it is sufficient to see that the matrices

\[
U^g_{z_1, x}(z^{1/2}, \theta, x_1) = \begin{bmatrix}
\cdots & \frac{z_1 \sin \theta}{z_1} \\
\cos \frac{\theta}{z_1} & z_1 \sin \frac{\theta}{z_1} \\
-z_1^{-1} \sin \frac{\theta}{z_1} & \cos \frac{\theta}{z_1} \\
\cdots & \frac{z_1 \sin \theta}{z_1}
\end{bmatrix},
\]

extend properly, which is clear. Hence, from now on we define $U^g_{z_1, x}(a, \theta, x_1)$ for all $\theta$ using the canonical continuous extension.

4.2. We can similarly define

\[
U^{X,g}_{z_1, x}(a, \theta, x_1) = D^{X,g}_{z_1}(\theta, x_1) U^g_{z_1, x}(a) D^{X,g}_{z_1}(\theta, x_1)^{-1}.
\]

Then the discussion is completely analogous. For example, for $X = R$ it yields

\[
U^R_{z_1, x}(z^n, \theta, x_1) =
\begin{bmatrix}
\cdots & 1 \\
1 - tz_1 & \cdots & \frac{t^n - z_1^n}{z_1} \\
-tz_1^{-1} & s^2 & \cdots & \frac{t^n - z_1^{-1} s^2}{s_1} \\
\cdots & \cdots & \cdots & \cdots \\
-tz_1^{-1} & s^2 & \cdots & \frac{t^2 s^2 z_1}{s_1} \\
\end{bmatrix};
\]

\[
U^R_{z_1, x}(z^{1/2}, \theta, x_1) =
\begin{bmatrix}
\cdots & 1 \\
1 - z_1^{-1} t & 1 \\
\cdots & \cdots \\
\end{bmatrix};
\]

etc. Note that $s^2 = 1 - (tz_1)(tz_1^{-1})$. This implies that in the cases $X = R, L$ the entries of matrices $U^{X,g}_{z_1, x}(a, \theta, x_1)$ are (infinite) expressions of $tz_1$, $t z_1^{-1}$ and the entries of $U^g_{z_1, x}(a)$ only. Similarly, in the cases $X = R, L$ the entries of matrices $U^{X,g}_{z_1, x}(a, \theta, x_1)$ are expressions of $\tau z_1$, $\tau z_1^{-1}$ and the entries of $U^g_{z_1, x}(a)$ only. The various deformations are quite similar: Due to

\[
U^{Y,g}_{z_1, x}(a, \theta, x_1) = \tilde{A}^{Y,g}_{z_1}(\theta, x_1) U^{X,g}_{z_1, x}(a, \theta, x_1) \tilde{A}^{X,g}_{z_1}(\theta, x_1)
\]

(for $\theta \notin \pi \mathbb{Z}$) and continuity we see that the matrices corresponding to various choices of $X$ differ from each other only in the environment of the coordinate axes.
4.3. **Lemma.** For \( n \geq 1, n \in \mathbb{N} \), it yields

\[
U^X_n(z^n, \theta, x_1) = A^X_n(\theta, x_1) \ldots A^X_{-1}(\theta, x_1) Z^n
\]

\[
= Z^n A^X_{n+\frac{1}{2}}(\theta, x_1) \ldots A^X_{-\frac{1}{2}}(\theta, x_1),
\]

\[
U^{\mathcal{A}}_{n}(z^{n-\frac{1}{2}}, \theta, x_1) = A^\mathcal{A}_n(\theta, x_1) \ldots A^\mathcal{A}_{-\frac{1}{2}}(\theta, x_1) Z^{n-\frac{1}{2}} =
\]

\[
= Z^{n-\frac{1}{2}} A^\mathcal{A}_{n-\frac{1}{2}}(\theta, x_1) \ldots A^\mathcal{A}_{-\frac{1}{2}}(\theta, x_1),
\]

\[
U^{\mathcal{A}}_{n}(z^{n-\frac{1}{2}}, \theta, x_1) = A^\mathcal{A}_{0}(\theta, x_1) \ldots A^\mathcal{A}_{n-1}(\theta, x_1) Z^{n-\frac{1}{2}} =
\]

\[
= Z^{n-\frac{1}{2}} A^\mathcal{A}_{n-1}(\theta, x_1) \ldots A^\mathcal{A}_{0}(\theta, x_1),
\]

\[
U^{\mathcal{A}}_{\mathcal{A}}(z^n, \theta, x_1) = A^{\mathcal{A}}_{0}(\theta, x_1) \ldots A^{\mathcal{A}}_{n-1}(\theta, x_1) Z^n = A^{\mathcal{A}}_{n}(\theta, x_1).
\]

**Proof.** It is sufficient to compute the case \( n = 1 \); the rest follows from the identities

\[
Z^n A^X_{k+n}(\theta, x_1) Z^{-n} = A^X_{k+n}(\theta, x_1),
\]

\[
Z^n A^{\mathcal{A}}_{k+n}(\theta, x_1) Z^{-n} = A^{\mathcal{A}}_{k+n}(\theta, x_1).
\]

\[\square\]

4.4. It is quite obvious that the deformations \( U^{X,g}_{n}(z_{i},z_{j})(a, \theta, x_1) \) can be constructed “by hand”, without the help of the matrices \( D^X_{n}(\theta, x_1) \). In fact, this is sort of necessary. In general, the deformations \( U^{R,g}_{n}(z_{i},z_{j})(a, \theta, x_1) \) and \( U^{L,g}_{n}(z_{i},z_{j})(a, \theta, x_1) \) do not respect unitarity but they are of rather algebraic nature, hence general. (They allow coefficients from more general rings.) Their behaviour quite differs from the original case for \( \theta = \frac{\pi}{2} \). The deformations \( U^{R,g}_{n}(z_{i},z_{j})(a, \theta, x_1) \) and \( U^{L,g}_{n}(z_{i},z_{j})(a, \theta, x_1) \) are sort of in the middle. Still, they are probably closer to the plain case \( X = \emptyset \).

4.5. We see that the case \( \theta = 0 \) means no deformation at all:

\[
U^{X,g}_{n}(z_{i},z_{j})(a, 0, x_1) = U^{g}_{z_{i},z_{j}}(a).
\]

The other interesting value is \( \theta = \frac{\pi}{2} \), where very special kind of degeneracies occur:

4.6. **Lemma.** It yields

\[
U^{g}_{z_{i},z_{j}}(a, \frac{\pi}{2}, x_1) = L^{\frac{\pi}{2}}_{z_{i}} \left\{ U^{g}_{z_{i},z_{j}}(a) \right\}^{g}_{z_{i},z_{j}} a(x_1) = L^{\frac{\pi}{2}}_{z_{j}} a(x_1),
\]

\[
U^{R,g}_{z_{i},z_{j}}(a, \frac{\pi}{2}, x_1) = L^{\frac{\pi}{2}}_{z_{i}} \left\{ U^{g}_{z_{i},z_{j}}(a) \right\}^{R,g}_{z_{i},z_{j}} a(x_1) = L^{\frac{\pi}{2}}_{z_{j}} a(x_1),
\]

\[
U^{L,g}_{z_{i},z_{j}}(a, \frac{\pi}{2}, x_1) = L^{\frac{\pi}{2}}_{z_{i}} \left\{ U^{g}_{z_{i},z_{j}}(a) \right\}^{L,g}_{z_{i},z_{j}} a(x_1) = L^{\frac{\pi}{2}}_{z_{j}} a(x_1),
\]

\[
U^{R,g}_{z_{i},z_{j}}(a, \frac{\pi}{2}, x_1) = L^{\frac{\pi}{2}}_{z_{i}} \left\{ U^{g}_{z_{i},z_{j}}(a) \right\}^{R,g}_{z_{i},z_{j}} S(\Delta_{z_{j},z_{i}}^{2}) a(x_1) = L^{\frac{\pi}{2}}_{z_{j}} a(x_1).
\]
\[
U_{z_i, z_j}^{L,g}(a, \pi, x_1) = \left. L_{z_i}^{-\frac{\pi}{2}} \right|_{z_i} \left( \left. U_{z_i', z_j'}^{g}(a) \right|_{z_i'} \cdot S(\Delta_{z_i', z_j'}^{x_1}(a)) a(x_1) \right)_{z_i, z_j}^{L,g} \cdot \left. L_{z_i}^{-\frac{\pi}{2}} \right|_{z_i}.
\]

4.7. For a multiplicatively invertible element \(a \in \mathfrak{B}^1\) we define

\[
O_{z_i}^{L,g}(a, \theta, x_1) = 1_{z_i} + t^2 L_{z_i}^{-\frac{\pi}{2}} \left( x_1 \mapsto L_{z_i}^{-\frac{\pi}{2}} \left( \begin{array}{cc}
0 & S_{z_i'}^g(\Delta_{z_i', z_i}^x(a(x_1))) \\
0 & 0
\end{array} \right) \right) L_{z_i}^{-\frac{\pi}{2}}.
\]

\[
O_{z_i}^{R,g}(a, \theta, x_1) = 1_{z_i} + t^2 L_{z_i}^{-\frac{\pi}{2}} \left( x_1 \mapsto L_{z_i}^{-\frac{\pi}{2}} \left( \begin{array}{cc}
0 & \pi_{z_i'}^R(\Delta_{z_i', z_i}^R(a(x_1))) \\
0 & 0
\end{array} \right) \right) L_{z_i}^{-\frac{\pi}{2}}.
\]

where the Poisson transform acts in the variable \(x_1\), entrywise in the matrices.

These are invertible matrices, and for \(X = L, R\) we set

\[
U_{z_i, z_j}^{X,g}(a, \theta, x_1) = O_{z_i}^{X,g}(a, \theta, x_1)^{-1} U_{z_i, z_j}^{X,g}(a, \theta, x_1).
\]

Then

a.) \(U_{z_i, z_j}^{X,g}(a, \theta, x_1)\) is a smooth perturbation of \(U_{z_i, z_j}^{X,g}(a, 0, x_1) = U_{z_i, z_j}^{g}(a)\).

b.) In the special case \(\theta = \frac{\pi}{2}\) it yields

\[
U_{z_i, z_j}^{L,g}(a, \frac{\pi}{2}, x_1) = L_{z_i}^{-\frac{\pi}{2}} \left\{ \begin{array}{c}
1_{z_i'} \\
S_{z_i'}^g(\Delta_{z_i', z_i}^x(a(x_1)^{-1})) \\
1
\end{array} \right\}_{z_i, z_j}^{L,g}.
\]

\[
= L_{z_i}^{-\frac{\pi}{2}} \left\{ \begin{array}{c}
1_{z_i'} \\
S_{z_i'}^g(\Delta_{z_i', z_i}^x(a(x_1)^{-1})) \\
1
\end{array} \right\}_{z_i, z_j}^{L,g}.
\]

And

\[
U_{z_i, z_j}^{R,g}(a, \frac{\pi}{2}, x_1) = L_{z_i}^{-\frac{\pi}{2}} \left\{ \begin{array}{c}
1_{z_i'} \\
S_{z_i'}^g(\Delta_{z_i', z_i}^x(a(x_1)^{-1})) \\
1
\end{array} \right\}_{z_i, z_j}^{R,g}.
\]

\[
= L_{z_i}^{-\frac{\pi}{2}} \left\{ \begin{array}{c}
1_{z_i'} \\
S_{z_i'}^g(\Delta_{z_i', z_i}^x(a(x_1)^{-1})) \\
1
\end{array} \right\}_{z_i, z_j}^{R,g}.
\]

c.) The entries of \(U_{z_i, z_j}^{L,g}(a, \theta, x_1)\) are infinite algebraic expressions of \(tz_1, tz_1^{-1}\) and the entries of \(U_{z_i, z_j}^{g}(a)\) and \(U_{z_i, z_j}^{g}(a)^{-1}\). (This is not entirely trivial to check, but it follows from Lemma 2.11 and Lemma 2.10)

d.) We remark that the correction terms \(O_{z_i}^{X,g}(a, \theta, x_1)\) are equal to \(1_{z_i}\) if \(a\) is a constant or to an \(\mathfrak{A}\)-linear combination of \(z^{1/2}\) and \(z^{-1/2}\). The matrices \(U_{z_i, z_j}^{X,g}(a, \theta, x_1)\) are not multiplicative. They are defined only to yield a better values for \(\theta = \frac{\pi}{2}\) than the “unprimed” versions.
In what follows we use the notation $\overline{X}$ for upper left indices when the “primed” variations $\iota L$ and $\iota R$ are also allowed along the ordinary terms $\iota, \overline{\iota}, \emptyset, \overline{\emptyset}, \overline{\iota}, \overline{\emptyset}$.

4.8. In general, suppose that we have a matrix operation which, given $a \in \mathfrak{B}^l$, $\theta$, $x_1$ it produces a $\mathbb{Z}_i \times \mathbb{Z}_j$ matrix

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, x_1).$$

Such a construction can also be written down by the algebraic kernel functions

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, x_1; y_1, y_2).$$

It may or may not have certain symmetries. Here are some of the possible symmetries described in the language of the matrices and the equivalent transcription in terms of transformation kernels:

a.) Rotation symmetry:

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(R^j a, \theta, x_1) = R_{\mathbb{Z}_i}^a \mathbb{X}_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, x_1 - \beta) R_{\mathbb{Z}_j}^a \mathbb{X}_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, x_1 - \beta; y_1 - \beta, y_2 - \beta).$$

b.) Reflection symmetry:

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(C a, \theta, x_1) = T_{\mathbb{Z}_i} \mathbb{Y}_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, -x_1) T_{\mathbb{Z}_j};$$

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(C a, \theta, x_1; y_1, y_2) = M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, -x_1; -y_1, -y_2).$$

c.) Central symmetry:

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, x_1) = M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, -\theta, x_1 + \pi);$$

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, \theta, x_1; y_1, y_2) = M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, -\theta, x_1 + \pi; y_1, y_2).$$

d.) Central degeneracy:

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, 0, x_1 - \beta)$$

does not depend on $\beta$;

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^a(a, 0, x_1 - \beta; y_1, y_2)$$

does not depend on $\beta$.

(The last two conditions together mean that the associated matrix depends essentially on $\iota z_1$, $\iota z_1^{-1}$ and $a$.)

e.) Duality symmetry: If the associated matrices come in a variety depending on $X$ then a possible duality symmetry is

$$M_{\mathbb{Z}_i, \mathbb{Z}_j}^{X, g, \text{opp}}(a, \theta, x_1) = M_{\mathbb{Z}_i, \mathbb{Z}_j}^{X, g, \text{opp}}(a, \theta, x_1) \|^\text{opp}.$$
4.9. For \( a \in \mathcal{B}^i \) we define
\[
E_{Z_i}^g(a, x_1) = \left\{ \begin{array}{l} \frac{g}{Z_i} \rightarrow x_1 \\ a(x_1) \end{array} \right\}^{g}_{Z_i}, \quad E_{Z_i}^{L,g}(a, x_1) = \left\{ \begin{array}{l} \frac{g}{Z_i} \rightarrow x_1 \\ a(x) \end{array} \right\}^{g,L}_{Z_i},
\]
and
\[
\hat{U}_{Z_i, a}^g(a) = \left\{ \begin{array}{l} U_{Z_i}^g(a) \\ 1 \end{array} \right\}_{Z_i}, \quad \hat{U}_{Z_i, a}^{L,g}(a) = \left\{ \begin{array}{l} U_{Z_i}^g(a) \\ 1 \end{array} \right\}_{Z_i}^{L,g}.
\]
One can make analogous definitions with "R". Then
\[
\hat{U}_{Z_i, a}^g(a, \frac{t}{2}, x_1) = E_{Z_i}^g(a, x_1) \hat{U}_{Z_i, a}^g(a) L_{\frac{t}{2}}^{g-\frac{t}{2}}.
\]
Furthermore, one can see that the matrices
\[
E_{Z_i}^g(a, x_1) L_{\frac{t}{2}}^{g-\frac{t}{2}} \quad \text{and} \quad L_{\frac{t}{2}}^{g-\frac{t}{2}} \hat{U}_{Z_i, a}^g(a) L_{\frac{t}{2}}^{g-\frac{t}{2}}
\]
have naturality, and the rotation, reflection and duality symmetries.

5. Kernel homotopies of rapidly decreasing matrices

5.1. Lemma. For rapidly decreasing \( \mathbb{Z}_i \times \mathbb{Z}_j \) matrices \( A^g \) the map
\[
P_{Z_i, Z_j}^{X,g}(\theta, x_1, x_2) : A^g \rightarrow P_{Z_i}^{X,g}(\theta, x_2)^{-1} A^g P_{Z_j}^{X,g}(\theta, x_1)^{-1} \cos^2 \theta
\]
extends smoothly to general \( \theta \).

Proof. a) Let us consider the case \( X = \emptyset, i = j = 0 \). For the elementary matrix \( e_{n,m} \) (\( n, m \geq 0 \)) it yields
\[
e_{n,m} \mapsto \begin{bmatrix} \cdots & 0 \\ \vdots & \vdots \end{bmatrix} \begin{bmatrix} z_1^{n-1} z_2^{-m} & z_1^{n-1} z_2^{-m+1} st & \cdots & z_1^{n-1} z_2^{-m+1} \frac{st^m}{z_1} \\ z_1^{n} z_2^{-m} st & z_1^{n-1} z_2^{-m+1} s^2 t^2 & \cdots & z_1^{n-1} z_2^{-m+1} s^2 t^{m+1} \\ \vdots & \vdots & \ddots & \vdots \\ z_1^{n} z_2^{-m} st^{n-1} & z_1^{n-1} z_2^{-m+1} s^2 t^n & \cdots & z_1^{n-1} z_2^{-m+1} s^2 t^{n+m-1} \\ z_1^0 z_2^{-m} s t^n & z_1^0 z_2^{-m+1} s^2 t^{n+1} & \cdots & z_1^0 z_2^{-m+1} s^2 t^{n+m} \\ & & \ddots & \vdots \\ & & & \ddots \end{bmatrix} \begin{bmatrix} \cdots \\ 0 \end{bmatrix}
\]
where \( t = \sin \theta \), \( s = \cos \theta \). The behaviour is similar for \( e_{n,m} \) from other quadrants. The result for \( a_{n,m} e_{n,m} \) is slightly more complicated: \( a_{n,m} \) should be inserted between \( z_1^n \) and \( z_2^m \). The other cases are all similar. \( \square \)
5.2. The case $\theta = 0$ means no deformation:

$$P_{X,\theta}^g(0, x_1, x_2)(A^g) = A^g.$$

5.3. **Lemma.** If the rapidly decreasing matrix $A^g$ has kernel function $A(x_1, x_2)$ then

$$P_{Z_i, Z_j}^g\left(\frac{x_i}{Z_i}, x_1, x_2\right)(A^g) = L_{Z_i}^{g, \frac{x_i}{Z_i}} \left\{ 0\right\}_{Z_i, Z_j}^{g, -\frac{x_i}{Z_i}} A(x_1, x_2)$$

and for $a \in B^1$ we define

$$K_{Z_i}^{X, g}(a, \theta, x_1, x_2) = U_{Z_i, Z_i}^{X, g}(a, \theta, x_1)L_{Z_i}^{g, \epsilon_1, x_2}(U_{Z_i, Z_i}^{X, g}(a, \theta, x_2))^{-1},$$

and for $a \in B^0$ we define

$$K_{Z_i}^{X, g}(a, \theta, x_1, x_2) = U_{Z_i, Z_i}^{X, g}(a, \theta, x_1)L_{Z_i}^{g, \epsilon_1, x_2}(U_{Z_i, Z_i}^{X, g}(a, \theta, x_2))^{-1}.$$

6. **Lemma.** It yields

$$K_{Z_i}^{X, g}(a, \theta, x_1, x_2) = 1_{Z_i} \cos \frac{x_1 - x_2}{2} + \left( H_{Z_i}^g + P_{Z_i}^X(\theta, x_1, x_2) J_{Z_i}^g(a) - H_{Z_i}^g \right) \sin \frac{x_1 - x_2}{2}.$$

**Proof.** First we prove the statement for $\theta \notin \pi \mathbb{A}$. From the definition of $U_{Z_i, Z_j}^{X, g}(a, \theta, x_1)$, and real multiplication commutativity with $b = \omega(x - x_1, \theta)$ we see that

$$K_{Z_i}^{X, g}(a, \theta, x_1, x_2) = F_{Z_i}^{X, g}(\theta, x_1)^{-1} U_{Z_i, Z_i}^{X, g}(a) F_{Z_i}^{X, g}(\theta, x_1) \left( 1_{Z_i} \cos \frac{x_1 - x_2}{2} + 

+ H_{Z_i}^g \sin \frac{x_1 - x_2}{2} \right) F_{Z_i}^{X, g}(\theta, x_2) U_{Z_i, Z_i}^{X, g}(a)^{-1} F_{Z_i}^{X, g}(\theta, x_2)^{-1}.$$

Applying Lemma 3.11 in order expand the three terms in the middle, applying real multiplication commutativity with $b = \cos \frac{x_1 + x_2}{2}$, and applying Lemma 3.11 again in order to write the terms back, we find

$$K_{Z_i}^{X, g}(a, \theta, x_1, x_2) = 1_{Z_i} \cos \frac{x_1 - x_2}{2} + H_{Z_i}^g \sin \frac{x_1 - x_2}{2} + 

+ F_{Z_i}^{X, g}(\theta, x_1)^{-1} \left( J_{Z_i}^g(a) - H_{Z_i}^g \right) F_{Z_i}^{X, g}(\theta, x_1)^{-1} \cos^2 \theta \sin \frac{x_1 - x_2}{2}.$$  

This proves the statement for $\theta \notin \pi \mathbb{A}$. Then, according to Lemma 5.1, the result extends by continuity. 

□
6.3. Theorem. $K^g_Z$ satisfies the properties of a good linearizing cocycle:

i.) The cocycle property:

$$K^g_Z(a, \theta, x_1, x_1) = 1_Z,$$

$$K^g_Z(a, \theta, x_1, x_2) = K^g_Z(a, \theta, x_1, x_3)K^g_Z(a, \theta, x_3, x_2).$$

ii.) $K^g_Z(a, \theta, x_1, x_2)$ is a smooth perturbation of

$$1_Z \cos \frac{x_1 - x_2}{2} + H^g_Z \sin \frac{x_1 - x_2}{2}.$$

iii.) Values at $\theta = 0$:

$$K^g_Z(a, 0, x_1, x_2) = 1_Z \cos \frac{x_1 - x_2}{2} + J^g_Z(a) \sin \frac{x_1 - x_2}{2}.$$

iv.) Values at $\theta = \frac{\pi}{2}$:

$$K^g_Z(d, \frac{\pi}{2}, x_1, x_2) = a(x_1) a(x_2)^{-1} e_{0,0} + 1_{Z\backslash \{0\}} \cos \frac{x_1 - x_2}{2} + H^g_Z \sin \frac{x_1 - x_2}{2}.$$

v.) If $G$ is a skew-involution then

$$K^g_Z(\cos \frac{\pi}{2} + G \sin \frac{\pi}{2}, x_1, x_2) = 1_Z \cos \frac{x_1 - x_2}{2} + H^g_Z[G] \sin \frac{x_1 - x_2}{2}.$$  

Proof. Points (i–iv) must be clear from the behaviour of $U^g_{Z_i, Z_j}(a, \theta, x_1)$ discussed previously. Points (ii–v) follow from Lemma 6.2 and the properties of the kernel homotopies. \hfill \Box

6.4. For $K^{X, g}_Z(a, \theta, x_1, x_2)$, in general, the very same statement can be made except in point (iv) where in general

$$K^{X, g}_Z(a, \frac{\pi}{2}, x_1, x_2) = E^{X, g}_Z(a, x_1) E^{X, g}_Z(a, x_2)^{-1}.$$  

Actually, this makes difference only for

$$K^{L, g}_Z(a, \frac{\pi}{2}, x_1, x_2) = \left\{ L^{g, \frac{x_1 - x_2}{2}}_Z L^{g, \frac{x_1}{2}}_Z S^{g}_Z(\Delta^{x_1 \theta}_Z a(x) - \Delta^{x_2}_Z a(x)) a(x_2)^{-1} \right\}^{X, g}_Z,$$

and for $K^{R, g}_Z(a, \frac{\pi}{2}, x_1, x_2)$, where a similar formula holds. Nevertheless, these are only very mild deformations of $K^g_Z(a, \frac{\pi}{2}, x_1, x_2)$; the difference is only a mostly inconsequential off-diagonal term.

There is an analogous statement for $K^{X, g}_Z(a, \theta, x_1, x_2)$, in general, but it is less nice in (iv), trivial in (v), and its intuitive meaning is not clear. In particular, it does not provide a naive linearization procedure for periodic loops.

6.5. It is worthwhile to examine the invariance properties of the linearizing cocycles constructed above. Compared to the discussion in [13], the difference is that we have two parameter variables: $x_1$ and $x_2$. The generalizations of naturality, the rotation, reflection, central symmetries and central degeneracy are straightforward, it is simply required to double the variable $x_1$; Along every expression $f(x_1)$ an other one, $f(x_2)$, must be inserted. However

e'.) Duality symmetry is better to be formulated, in general, as

$$M^{X, g}_{Z_i, Z_j}(a^{opp}, \theta, x_1, x_2) = M^{X, g}_{Z_j, Z_i}(a^{-1}, \theta, x_2, x_1)^{opp},$$

$$M_{Z_i, Z_j}(a^{opp}, \theta, x_1, x_2, y_1, y_2) = M_{Z_j, Z_i}(a^{-1}, \theta, x_2, x_1, y_2, y_1)^{opp}.$$  

in terms of matrices and kernels respectively.

Then $L^{g, \frac{x_1 - x_2}{2}}_Z$, and, more generally, $K^{X, g}_Z(a, \theta, x_1, x_2)$ have all these symmetries.
6.6. Even at this point it becomes apparent that it might be useful to formulate the linearization procedure entirely in terms of cocycles. For a loop $a$ we might form the trivial cocycle

$$(x, \tilde{x}) \mapsto \tilde{a}(x, \tilde{x}) = a(x)a(\tilde{x})^{-1}.$$ 

If $(x, \tilde{x}) \to c(x, \tilde{x})$ is a cocycle then we define its dual cocycle by

$$(x, \tilde{x}) \mapsto c^\text{dual}(x, \tilde{x}) = c(\tilde{x}, x)^\text{opp}.$$ 

$e^n$.) If we have a matrix construction $M^X_{\beta,\gamma}(x, \theta, x_1, x_2)$ depending on a cocycle $c$ then we define duality invariance by

$$M^X_{\beta,\gamma}(c^\text{dual}, \theta, x_1, x_2) = M^X_{\beta,\gamma}(c, \theta, x_2, x_1)^\text{opp}.$$ 

in terms of matrices and kernels respectively.

First, it is clear from the real representation that $K_{\beta,\gamma}^X(a, \theta, x_1, x_2)$ depends only on the cocycle $\tilde{a}$ and not on its trivialization $a$. Hence, we could have written $K_{\beta,\gamma}^X(\tilde{a}, \theta, x_1, x_2)$. Then we may observe that it satisfies duality invariance in the above sense. In fact, we can formulate rotated cocycles by $R^\beta c(x_1, x_2) = c(x_1 - \beta, x_2 - \beta)$, etc., and we could have formulated rotation invariance, etc., in these terms.

If we consider linearization of loops in strict sense then this corresponds to finding a trivialization of the linearizing cocycle. If we choose such a trivialization then we generally lose some kind of invariance property.

6.7. When we consider loops as pointed maps, i.e. when $a(0) = 1$ holds then

$$\tilde{K}^g_Z(a, \theta, x_1) = K^g_Z(a, \theta, x_1, 0)$$

yields a linearizing homotopy which leaves that class of loops invariant. It yields a trivialization of the cocycle in the sense that

$$K^g_Z(a, \theta, x_1, x_2) = \tilde{K}^g_Z(a, \theta, x_1)K^g_Z(a, \theta, x_2)^{-1}.$$ 

Here both the trivial loop extension

$$\tilde{K}^g_Z(a, \frac{x}{2}, x_1) = a(x_1)e_{0,0} + 1_{Z\setminus\{0\}} \cos \frac{x_1}{2} + H_Z^g \sin \frac{x_1}{2}$$

and the linearized loop

$$\tilde{K}^g_Z(a, 0, x_1) = 1_Z \cos \frac{x_1}{2} + J_Z^g(a) \sin \frac{x_1}{2}$$

are nice. Its nice properties are immediate from the statement of Theorem 6.6. However, there is no apparent rotation invariance in this construction. There are natural variants $\tilde{K}^X_{\beta,\gamma}(a, \theta, x_1)$ of this construction, of course.

6.8. We have another construction which is less invariant than $K^g_Z(a, \theta, x_1, x_2)$ but in many ways better than $\tilde{K}^g_Z(a, \theta, x_1)$:

Let us consider

$$N^g_Z(a, \theta, x_1) = U_{Z,\mathcal{X}}^g(a, \theta, x_1)U_{Z,\mathcal{X}}^g(a)J_{Z,\mathcal{X}}^g(a \frac{x_1}{2}, x_1)^{-1}U_{Z,\mathcal{X}}^g(a, \theta, x_1)^{-1}E_Z^g(a, x_1).$$

It has the trivializing property

$$K^g_Z(a, \theta, x_1, x_2) = N^g_Z(a, \theta, x_1)N^g_Z(a, \theta, x_2)^{-1}.$$ 

The advantage of this function is that it directly linearizes

$$N^g_Z(a, \frac{x}{2}, x_1) = a(x_1)e_{0,0} + 1_{Z\setminus\{0\}} \cos \frac{x_1}{2} + H_Z^g \sin \frac{x_1}{2},$$

and the linearized loop

$$N^g_Z(a, 0, x_1) = 1_Z \cos \frac{x_1}{2} + J_Z^g(a) \sin \frac{x_1}{2}$$

are nice. Its nice properties are immediate from the statement of Theorem 6.6. However, there is no apparent rotation invariance in this construction. There are natural variants $N^X_{\beta,\gamma}(a, \theta, x_1)$ of this construction, of course.
although the linearized loop
\[ N_0^g(a, 0, x_1) = U_{z,x}^g(a) \left( 1_z \cos \frac{x_1}{2} + H_z^g \sin \frac{x_1}{2} \right) \tilde{U}_{z,x}^g(a)^{-1} \]
is not particularly nice in itself. Nevertheless, the failing of the construction to be rotation invariant is not worse than that of \( \hat{X} = \hat{X} \) itself: the equation of rotation invariance holds up to multiplication by \( L_A^g \). What we can say that here linearization “acts” on the left side while the right side is undetermined up to multiplication by a constant. One can consider variants \( N_X^g(a, \theta, x_1) \) which may or may not have better properties (like algebraic parametrizability in terms of \( t, \) etc.).

We remark that in the special cases \( a = z^{\pm 1/2} \) it yields
\[
N_X^g(z^{1/2}, \theta, x_1) = L_{z}^g[z^{1/2}] A_{z}^X(\theta) A_{z}^X(\frac{\theta}{2})^{-1},
\]
\[
N_X^g(z^{-1/2}, \theta, x_1) = L_{z}^g[z^{-1/2}] A_{z}^X(-\theta) A_{z}^X(-\frac{\theta}{2})^{-1}.
\]
(The primed versions do not differ in these cases.)

6.9. Of the various constructions above the cases \( X = R, L, tR, tL \) allow easy algebraization, ie. more general coefficient rings. In these cases \( U_{z_1}^\theta(a, \theta, x_1) \), \( E_X^\theta(a, x_1) \) and \( U_{z_1}^\theta(a) \) are expressed in terms of \( tz_1, tz_1^{-1} \), and the entries of \( U_{z_1}^\theta(a) \); not even division by 2 is required.

7. Non-commutative cyclic loops

Having gained some experience with linearizing loops we can try to generalize the results above. The main technique we use is simply suppression of variables, which allows to formulate statements in greater generality.

7.1. Consider a superalgebra \( \mathfrak{A} = \mathfrak{A}^+ \oplus \mathfrak{A}^- \), and consider an extension \( \mathfrak{A}[z, z^{-1}] \) where that \( z \) is a cyclic variable which does not necessarily commute with elements of \( \mathfrak{A} \) but \( z \mathfrak{A}^+ z^{-1} \subset \mathfrak{A}^+ \), \( z^{-1} \mathfrak{A}^+ z \subset \mathfrak{A}^+ \) and \( z \mathfrak{A}^- \subset \mathfrak{A}^- \), \( z^{-1} \mathfrak{A}^- z^{-1} \subset \mathfrak{A}^- \). Then \( \mathfrak{A}[z, z^{-1}] \) is naturally \( \mathbb{Z}_2 \times \mathbb{Z} \) graded, and assume that it is endowed by a (not necessary complete) locally convex algebra topology and \( \mathfrak{B}^{0,0} \) is the algebra of rapidly decreasing sequences regarding the grading. Ie. what we consider smooth loops with a non-commuting loop variable. One wonders if some kind of linearization is possible in this case.

7.2. We can introduce the symbol \( z^{1/2} \), and we can take \( \mathfrak{B}^0 = \mathfrak{B}^{0,0} \oplus \mathfrak{B}^{1,1} = \mathfrak{A}[z, z^{-1}] \oplus z^{1/2} \mathfrak{A}[z, z^{-1}] z^{1/2} \) and \( \mathfrak{B}^1 = \mathfrak{B}^{0,1} \oplus \mathfrak{B}^{1,0} = \mathfrak{A}[z, z^{-1}] z^{1/2} \oplus z^{1/2} \mathfrak{A}[z, z^{-1}]. \)

Regarding linearizations, what we have to do is to provide analogous formulas with a suppression of variables. That means we have to replace “\( y, x_1 \)” by \( z \) and “\( g, \frac{\pi}{2} \)” by \( z^{1/2} \) as much as it possible. But we have to keep in mind that, say, \( \mathfrak{A}^+ \) is not invariant for conjugation by \( z^{1/2} \).

If \( a \in \mathfrak{B}^{1,j} \) then according to the choice of basis \( e_n \leftrightarrow z^n \) the representing matrix of multiplication by \( a \) on the left is a matrix \( U_{z_1}^x(a) \). Let \( a = \sum_{n \in \mathbb{Z}_1} z^n a_n \) where \( a_n \in \mathfrak{A} \) or \( z^{1/2} \mathfrak{A}^+ z^{-1/2} \oplus z^{1/2} \mathfrak{A}^- z^{-1/2} \). The effect of multiplication by \( a \) on the left is given as \( az^k = \sum_{n \in \mathbb{Z}_1} z^{n+k}(z^{-k} a_n^+ z^k) + z^{n-k}(z^k a_n z^k) \). According to the correspondence of
Symbolic rotation actions are given by $R_{Z_i}^Z = \sum_{n \in Z_i} z^{-n}e_{n,n}$.

We can proceed further by other definitions like $L_\mathcal{Z}[u] = z1_{Z^-} + u e_{0,0} + z^{-1}1_{Z^+}$, $L_\mathcal{X}^Z = z1_{\mathcal{X}^-} + z^{-1}1_{\mathcal{X}^+}$, etc. Proceeding analogously to the skew-complex case then we see that $U_{Z_i, Z_j}(a, \theta)$ will also be of rotationally Toeplitz type, etc. It may be strange that the expression $U_{Z_i, Z_j}(a, \theta)$ would correspond to $U_{Z_i, Z_j}(a, \theta, x)$, i.e. $x$ appears as a variable of $a(x)$ and also appears as an external parameter (as we have an apparent shortage of variables). This, in fact, causes no problems, the first occurrence of $x$ is closed relative to the second. It is similar to expressions of type $\exists x \mathcal{F}(x, x_1, \ldots, x_n; \exists x \mathcal{G}(x, x_1, \ldots, x_n))$ in logic.

Regarding the special case $\theta = \frac{\pi}{2}$ we cannot write expressions which are similarly compact (as we cannot use $g$, we cannot simply take difference functions, etc.), nevertheless the structure of the matrices remains highly special, analogous to the skew-complex case. Ultimately, taking direct analogues of the various matrices as $Z_i \times Z_j$ matrices (and not as block decompositions) we can define $\hat{U}_{Z_i, Z_j}(a, \theta)$ etc. In theory this requires that we expand all our matrices as in, say, point 1 and make sure that we can proceed. For $a \in \mathfrak{B}^{0,1}$ we find that $U_{Z_i, Z_j}(a, \theta)$ is matrix with coefficients from $\mathfrak{B}^{0,1}$ and $\hat{U}_{Z_i, Z_j}(a, \theta)$ is a matrix with coefficients from $z^{1/2}\mathfrak{A}^{+}z^{-1/2} \oplus z^{1/2}\mathfrak{A}^{-}z^{-1/2}$. Then we can indeed present a linearization:

7.4. For $a \in \mathfrak{B}^{0,1}$ we can define

$$N_{\mathfrak{B}}^Z(a, \theta) = U_{Z,\mathfrak{A}}^Z(a, \theta) L_{\mathfrak{A}}^{Z^{1/2}} \hat{U}_{Z,\mathfrak{A}}^Z(a, \theta)^{-1} = U_{Z,\mathfrak{A}}^Z(a, \theta) U_{Z,\mathfrak{A}}^Z(a, \frac{\pi}{2})^{-1} L_{\mathfrak{Z}}^{Z^{1/2}} [a(z)].$$

7.5. Theorem. It yields:

i.) $N_{\mathfrak{B}}^Z(a, \theta)$ is rapidly decreasing perturbation of $L_{\mathfrak{Z}}^{Z^{1/2}[0]}$ with entries from $\mathfrak{B}^{0,1}$.

ii.) $N_{\mathfrak{B}}^Z(a, 0) = U_{Z,\mathfrak{A}}^Z(a) L_{\mathfrak{A}}^{Z^{1/2}} \hat{U}_{Z,\mathfrak{A}}^Z(a)^{-1}$ has entries from $\mathfrak{Z}^{1/2} \oplus \mathfrak{Z}^{-1/2}$.

iii.) $N_{\mathfrak{B}}(a, \frac{\pi}{2}) = L_{\mathfrak{Z}}^{Z^{1/2}} [a(z)].$

iv.) In the special cases $a = z^{\pm1/2}$ it gives $N_{\mathfrak{B}}^Z(z^{1/2}, \theta) = L_{\mathfrak{Z}}^{Z^{1/2}} [z^{1/2}] \mathfrak{A}_1^Z(\theta) \mathfrak{A}_1^Z(\frac{\pi}{2})^{-1}$, and $N_{\mathfrak{B}}^Z(z^{-1/2}, \theta) = L_{\mathfrak{Z}}^{Z^{1/2}} [z^{-1/2}] \mathfrak{A}_1^Z(-\theta) \mathfrak{A}_1^Z(-\frac{\pi}{2})^{-1}$. □

Other variants of this linearization procedure exist naturally. We can linearize loops $a \in \mathfrak{B}^{0,0}$ in more conventional sense by taking, for example, $N_{\mathfrak{B}}^Z(az^{-1/2}, \theta)z^{1/2}$.
8. Linearization in terms of transformation kernels

8.1. Let us consider loops and matrix actions on them using Fourier expansions with respect to \( g = i \). Let \( \mathbf{a} = \sum_{k \in \mathbb{Z}_j} a_k \mathbf{e}_k \), \( \mathbf{C} = \sum_{n \in \mathbb{Z}_i, m \in \mathbb{Z}_j} c_{n,m} \mathbf{e}_{n,m} \) be rapidly decreasing matrices. Through the correspondences
\[
\mathbf{a} = \sum_{k \in \mathbb{Z}_j} a_k \mathbf{e}_k \leftrightarrow a(x) = \sum_{k \in \mathbb{Z}_j} a_k e^{ikx}
\]
\[
\mathbf{C} = \sum_{n \in \mathbb{Z}_i, m \in \mathbb{Z}_j} c_{n,m} \mathbf{e}_{n,m} \leftrightarrow C(x, \tilde{x}) = \sum_{n \in \mathbb{Z}_i, m \in \mathbb{Z}_j} c_{n,m} e^{i(nx-m\tilde{x})}
\]
we see that various matrix operations can be represented according to
\[
\mathbf{C} \mathbf{a} \leftrightarrow C \ast a(x) = \int_{y=0}^{2\pi} C(x, y) a(y) \frac{dy}{2\pi},
\]
\[
\mathbf{C}_1 \mathbf{C}_2 \leftrightarrow C_1 \ast C_2(x, \tilde{x}) = \int_{y=0}^{2\pi} C_1(x, y) C_2(y, \tilde{x}) \frac{dy}{2\pi},
\]
\[
\mathbf{C}^T \leftrightarrow C^T(x, \tilde{x}) = C(-\tilde{x}, -x).
\]
As transformation kernels do not depend on any choice like \( g = i \) we see that the operations "\( \ast \)" and "\( T \)" are entirely canonical. After this integration will be understood over \([0, 2\pi]\) (or over a compactly supported smooth distribution such that the sum of its translates by \( 2\pi \mathbb{Z} \) is the constant 1 function.)

Using these observations, in what follows, we will describe our linearizations in terms of transformation kernels. Instead of writing a transformation kernel as \( C(x_1, x_2) \) we will sometimes write it as \( C[x_1, \alpha] \) where \( \alpha = x_1 - x_2 \). Then
\[
C[x_1, \alpha] = C(x_1, x_1 - \alpha) = C(x_1, x_2) = C[x_1, x_1 - x_2].
\]

In this formalism
\[
C \ast a(x) = \int_\beta C(x, \beta) a(x - \beta) \frac{d\beta}{2\pi} = \int_\beta C[x, \beta] R^\beta a(x) \frac{d\beta}{2\pi},
\]
\[
C_1 \ast C_2[x, \alpha] = \int_\beta C_1[x, \beta] C_2[x - \beta, \alpha - \beta] \frac{d\beta}{2\pi} = \int_\beta C_1[x, \beta] R^\beta C_2[x, \alpha - \beta] \frac{d\beta}{2\pi},
\]
\[
C^T[x, \alpha] = C[\alpha - x, \alpha];
\]
here \( R^\alpha \) means rotated in the first variable.

8.2. Actually, the most natural operations cannot be represented by smooth transformation kernels. We can allow transformation kernels which are distributions in the variable \( \alpha \), and there is a smooth dependence in \( x \). For \( b \in \mathfrak{B}^j \) the standard multiplication and rotation actions are given by
\[
U_{\mathbb{Z}_i, \mathbb{Z}_j}(a) \ast b(x) = a(x) b(x) = \int_\alpha a(x) \delta^0_{\mathbb{Z}_j}(\alpha) R^\alpha b(x),
\]
\[
R^\beta_{\mathbb{Z}_j} \ast b(x) = \int_\alpha \delta^\beta_{\mathbb{Z}_j}(\alpha) R^\alpha b(x);
\]
some Dirac \( \delta \)-functions appear. Hence we find that the kernels are
\[
U_{\mathbb{Z}_i, \mathbb{Z}_j}(a)[x, \alpha] = a(x) \delta^0_{\mathbb{Z}_j}(\alpha) = U_{\mathbb{Z}_i, \mathbb{Z}_j}(a)(x, \tilde{x}) = a(x) \delta^0_{\mathbb{Z}_j}(x - \tilde{x}) = a(\tilde{x}) \delta^0_{\mathbb{Z}_j}(x - \tilde{x}),
\]
\[
R^\beta_{\mathbb{Z}_j}[x, \alpha] = \delta^\beta_{\mathbb{Z}_j}(\alpha) = R^\beta_{\mathbb{Z}_j}[x, \tilde{x}] = \delta^\beta_{\mathbb{Z}_j}(x - \tilde{x}).
\]
For a transformation kernel $C$ representing a $Z_i \times Z_j$ matrix one finds that the following identities are computationally useful:

\[
C \ast U_{Z_i, Z_k}(a)[x, \alpha] = C[x, \alpha]a(x - \alpha) = C \ast U_{Z_i, Z_k}(a)(x, \bar{x}) = C(x, \bar{x})a(\bar{x}),
\]

\[
U_{Z_i, Z_k}(a) \ast C[x, \alpha] = a(x)C[x, \alpha] = U_{Z_i, Z_k}(a) \ast C(x, \bar{x}) = a(x)C(x, \bar{x}),
\]

\[
R^\beta_{Z_i} \ast (R^\beta_{Z_j})[x, \alpha] = C[x - \beta, \alpha] = R^\beta_{Z_i} \ast (R^\beta_{Z_j})(x, \bar{x}) = C(x - \beta, \bar{x} - \beta).
\]

In particular, a transformation is rotation invariant if and only if its kernel depends only on $\alpha$.

### 8.3. Such transformations are, for example, the Poisson, Poisson-Hilbert, and Hilbert transformations.

Using the notation $w = e^{i\alpha}$ we find that their kernels are

\[
P^*_Z[x, \alpha] = \frac{1 - r^2}{(1 - rw)(1 - rw^{-1})} = \frac{1 - r^2}{1 - 2r \cos \alpha + r^2},
\]

\[
H^*_Z[x, \alpha] = \frac{-ir(w - w^{-1})}{(1 - rw)(1 - rw^{-1})} = \frac{2r \sin \alpha}{1 - 2r \cos \alpha + r^2}.
\]

\[
H_Z[x, \alpha] = PV \frac{w + 1}{w - 1} = PV \cot \frac{\alpha}{2}.
\]

\[
P^r_Z[x, \alpha] = \frac{r^{1/2}(1 - r)(w^{1/2} + w^{-1/2})}{(1 - rw)(1 - rw^{-1})} = \frac{2r^{1/2}(1 - r) \cos \frac{\alpha}{2}}{1 - 2r \cos \alpha + r^2},
\]

\[
H^r_Z[x, \alpha] = \frac{-ir^{1/2}(1 + r)(w^{1/2} - w^{-1/2})}{(1 - rw)(1 - rw^{-1})} = \frac{2r^{1/2}(1 + r) \sin \frac{\alpha}{2}}{1 - 2r \cos \alpha + r^2},
\]

\[
H_Z[x, \alpha] = PV \frac{2i}{w^{1/2} - w^{-1/2}} = PV \cosec \frac{\alpha}{2}.
\]

So, we can forget the first variable $x$. After this we omit the sign $PV$.

### 8.4. The Dirac delta and Cauchy principal values are among the very mildest singularities. Yet, heuristical computations with them should be taken with a grain of salt. For example, when one takes the square of the Hilbert transform one finds

\[
cot \frac{\beta}{2} \cot \frac{\alpha - \beta}{2} - 1 = \begin{cases} 
- \cosec^2 \frac{\beta}{2} & \text{if } \alpha \in 2\pi \mathbb{Z}, \\
\cot \frac{\alpha}{2} \left( \cot \frac{\beta}{2} + \cot \frac{\alpha - \beta}{2} \right) & \text{if } \alpha \notin 2\pi \mathbb{Z},
\end{cases}
\]

and

\[
\cosec \frac{\beta}{2} \cosec \frac{\alpha - \beta}{2} = \begin{cases} 
- \cosec^2 \frac{\beta}{2} & \text{if } \alpha \in 4\pi \mathbb{Z}, \\
\cosec \frac{\alpha}{2} \left( \cosec \frac{\beta}{2} + \cosec \frac{\alpha - \beta}{2} \right) & \text{if } \alpha \notin 2\pi \mathbb{Z}, \\
\cosec \frac{\beta}{2} & \text{if } \alpha \in 4\pi \mathbb{N},
\end{cases}
\]

from which

\[
\int \cot \frac{\beta}{2} \cot \frac{\alpha - \beta}{2} \frac{d\beta}{2\pi} - 1 = -\delta^0_Z(\alpha)
\]

and

\[
\int \cosec \frac{\beta}{2} \cosec \frac{\alpha - \beta}{2} \frac{d\beta}{2\pi} = -\delta^0_Z(\alpha).
\]

are quite believable but it needs some faith regarding the singular values.

If one wants to be more precise, it is better to rely on the following lemma in order to approximate properly:
8.5. Lemma. For $0 \leq r_2 < r_1$

\[ P_{Z_{r_1}}^{r_1}(\beta)P_{Z_{r_2}}^{r_2}(\alpha - \beta) = \]

\[ = \frac{P_{Z_{r_1}}^{r_1}(\beta) + P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha) + \frac{P_{Z_{r_2}}^{r_2}(\beta) - P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha) \]

\[- \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_2}}^{r_1}(\alpha) + \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha), \]

\[ P_{Z_{r_1}}^{r_1}(\beta)H_{Z_{r_2}}^{r_2}(\alpha - \beta) = \]

\[ = \frac{P_{Z_{r_1}}^{r_1}(\beta) + P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} H_{Z_{r_1}}^{r_1}(\alpha) + \frac{P_{Z_{r_2}}^{r_2}(\beta) - P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} H_{Z_{r_1}}^{r_1}(\alpha) \]

\[ + \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_2}}^{r_1}(\alpha) - \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha), \]

\[ H_{Z_{r_1}}^{r_1}(\beta)P_{Z_{r_2}}^{r_2}(\alpha - \beta) = \]

\[ = \frac{P_{Z_{r_1}}^{r_1}(\beta) + P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} H_{Z_{r_1}}^{r_1}(\alpha) + \frac{-P_{Z_{r_2}}^{r_2}(\beta) + P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} H_{Z_{r_1}}^{r_1}(\alpha) \]

\[ + \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_2}}^{r_1}(\alpha) + \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha), \]

\[ H_{Z_{r_1}}^{r_1}(\beta)H_{Z_{r_2}}^{r_2}(\alpha - \beta) - \delta_{t_0} \cdot 1 = \]

\[ = - \frac{P_{Z_{r_1}}^{r_1}(\beta) + P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha) + \frac{P_{Z_{r_2}}^{r_2}(\beta) - P_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} P_{Z_{r_1}}^{r_1}(\alpha) \]

\[ + \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} H_{Z_{r_1}}^{r_1}(\alpha) + \frac{H_{Z_{r_1}}^{r_1}(\beta) + H_{Z_{r_2}}^{r_2}(\alpha - \beta)}{2} H_{Z_{r_1}}^{r_1}(\alpha), \]

and similarly with $\beta$ and $\alpha - \beta$ interchanged. \□

8.6. If a transformation kernel also depends on an additional parameter $x_1$ then we can write it in alternative ways, as

\[ C(x_1, y_1, y_2) = C[x_1, y_1, \beta] = C[[x_1, \eta, \beta]], \]

where $\beta = y_1 - y_2, \eta = x_1 - y_1$. In this case, for the multiplication it yields

\[ C_1 * C_2(x_1, y_1, y_2) = \int_u C_1(x_1, y_1, u)C_2(x_1, u, y_2) \frac{du}{2\pi} = \]

\[ = C_1 * C_2[x_1, y_1, \beta] = \int_\chi C_1[x_1, y_1, \chi]C_2[x_1, y_1 - \chi, \beta + \chi] \frac{d\chi}{2\pi} = \]

\[ = C_1 * C_2[[x_1, \eta, \beta]] = \int_\chi C_1[[x_1, \eta, \chi]]C_2[[x_1, \eta + \chi, \beta + \chi]] \frac{d\chi}{2\pi}. \]

8.7. From the kernel of the Hilbert transform we find that for $a \in \mathcal{B}^i$

\[ J_{Z_{r_1}}(a)(y_1, y_2) = \frac{a(y_1)a(y_2)^{-1}}{\sin \frac{y_1 - y_2}{2}} = \frac{a(y_1, y_2)}{\sin \frac{y_1 - y_2}{2}}. \]

Let us concentrate to the case $X = \emptyset$. We find

\[ F_Z(\theta, x_1) = \delta_{Z}^0(y_1 - y_2)(1 - \cos(y_1 - x_1)\sin \theta) + \frac{\sin(y_1 + y_2 - x_1)}{\sin \frac{y_1 - y_2}{2}} \sin \theta + (\cos \theta - 1), \]
\[ F_Z^0(\theta, x_2) = \delta_Z^0(y_1 - y_2)(1 - \cos(y_2 - x_2)\sin \theta - \frac{\sin(y_1 + y_2 - x_2)}{\sin y_1 - y_2} \sin \theta + (\cos \theta - 1). \]

From Lemma 6.2 we know that
\[ K_Z^0(a, \theta, x_1, x_2) = \cos \frac{x_1 - x_2}{2} + \sin \frac{x_1 - x_2}{2} \left( H_Z^0 + \right. \]
\[ + F_Z^0(\theta, x_1) U_Z^0(\omega(x - x_1, \theta)^{-1})(F_Z(a) - H_Z^0) U_Z^0(\omega(x - x_2, \theta)^{-1}) F_Z^0(\theta, x_2) \cos^2 \theta \right) \]

8.8. Then we can write down the transformation kernel
\[ K_Z(a, \theta, x_1, x_2, y_1, y_2) = \cos \frac{x_1 - x_2}{2} \delta_Z^0(y_1 - y_2) + \sin \frac{x_1 - x_2}{2} \left( \cot \frac{y_1 - y_2}{2} + \right. \]
\[ + \left. \left[ 1 - \cos(y_1 - x_1) \sin \theta \left( \frac{\sin \frac{y_1 - y_2}{2}}{\omega(y_1 - x_1, \theta)} - \cos \frac{\sin \frac{y_1 - y_2}{2}}{\omega(y_1 - x_1, \theta)} \right) \right] \frac{1 - \cos(y_2 - x_2) \sin \theta}{\omega(y_2 - x_2, \theta)} + \right. \]
\[ + \int \frac{1}{u_1} \left( \frac{\sin \frac{y_1 + y_2}{2} - x_1}{\sin \frac{y_1 - y_2}{2}} \sin \theta + (\cos \theta - 1) \right) \left( \frac{\sin \frac{y_1 + y_2}{2}}{\omega(y_1 - x_1, \theta) \sin \frac{y_1 - y_2}{2}} \right) \frac{1 - \cos(y_2 - x_2) \sin \theta}{\omega(y_2 - x_2, \theta)} + \ldots \right] \cos^2 \theta \right) \]

It has the following properties:

i.) It is a smooth perturbation of
\[ \cos \frac{x_1 - x_2}{2} \delta_Z^0(y_1 - y_2) + \sin \frac{x_1 - x_2}{2} \cot \frac{y_1 - y_2}{2} \]

ii.) It has the cocycle property:
\[ K(a, \theta, x_1, y_1, y_2) = \delta_Z^0(y_1 - y_2), \]
\[ K_Z(a, \theta, x_1, x_2, y_1, y_2) = \int K(a, \theta, x_1, x_3, y_1, u) K(a, \theta, x_3, x_2, u, y_2). \]

iii.) Value at \( \theta = 0 \):
\[ K(a, 0, x_1, x_2, y_1, y_2) = \cos \frac{x_1 - x_2}{2} \delta_Z^0(y_1 - y_2) + \sin \frac{x_1 - x_2}{2} \cot \frac{y_1 - y_2}{2} \tilde{a}(y_1, y_2). \]

iv.) Value at \( \theta = \frac{\lambda}{2} \):
\[ K(a, \frac{\lambda}{2}, x_1, x_2, y_1, y_2) = \cos \frac{x_1 - x_2}{2} \left( \delta_Z^0(y_1 - y_2) - 1 \right) + \sin \frac{x_1 - x_2}{2} \cot \frac{y_1 - y_2}{2} + \tilde{a}(x_1, x_2). \]

v.) If \( G \) is a skew-involution then
\[ K(\cos \frac{\lambda}{2} + G \sin \frac{\lambda}{2}, \theta, x_1, y_1, y_2) = \cos \frac{x_1 - x_2}{2} \delta_Z^0(y_1 - y_2) + \sin \frac{x_1 - x_2}{2} \left( \cos \frac{y_1 - y_2}{2} + G \right) . \]

The various symmetries are manifested in the kernel. For example, rotation symmetry is related to the fact that the transformation can be expressed relatively simply by the cocycle \( \tilde{a}, \), and \( x_1, \) and \( \alpha = x_1 - x_2, \eta = x_1 - y_1, \beta = y_1 - y_2. \)
9. Linearization of $S^1$-cocycles

9.1. Fourier expansion of $S^1$-actions. Assume that $\mathcal{B} = \mathcal{B}^0 \oplus \mathcal{B}^1$ is a $\mathbb{Z}_2$-graded algebra. Assume that it allows a smooth one-parameter group of automorphisms $R : \mathbb{R} \times \mathcal{B} \to \mathcal{B}$ which is periodic on $\mathcal{B}^0$ and skew-periodic on $\mathcal{B}^1$. We use $R^\alpha$ to denote the action of $R$ belonging to the parameter value $\alpha$. Then the periodicity conditions are $R^{2\pi}b = b$ for $b \in \mathcal{B}^0$ and $R^{2\pi}b = -b$ for $b \in \mathcal{B}^1$.

This situation is characterized by the lack of an explicit loop variable $x$. Nevertheless, even in this case one can apply Fourier decomposition in terms of the rotation action. Indeed, due to smoothness the action of $R$ allows a Fourier decomposition

$$R^\alpha a = \frac{1}{2}a_{[0]} + \sum_{n \in \mathbb{Z}^+} a_{[n]} \cos n\alpha + H_Z a_{[n]} \sin n\alpha \quad \text{for} \quad a \in \mathcal{B}_0,$$

$$R^\alpha a = \sum_{n \in \mathbb{Z}^+} a_{[n]} \cos n\alpha + H_Z a_{[n]} \sin n\alpha \quad \text{for} \quad a \in \mathcal{B}_1;$$

where the Hilbert transforms are defined by

$$H_Z a = \int_\beta \cot \frac{\beta}{2} R^\alpha a \frac{d\beta}{2\pi} \quad \text{and} \quad H_Z a = \int_\beta \cosec \frac{\beta}{2} R^\alpha a \frac{d\beta}{2\pi};$$

and more particularly,

$$a_{[n]} = 2 \int_\beta \cos n\beta R^\alpha a \frac{d\beta}{2\pi} \quad \text{and} \quad H_Z a_{[n]} = 2 \int_\beta \sin n\beta R^\alpha a \frac{d\beta}{2\pi}.$$

(Cf.: If $b \in \mathcal{B}^0 = C_\mathbb{Z}(\mathcal{A})$ and

$$b(x) = b_0 + \sum_{n \in \mathbb{Z}^+} (b_n \cos nx + b_{-n} \sin nx)$$

then

$$R^\alpha b(x) = b_0 + \sum_{n \in \mathbb{Z}^+} (b_n \cos (x-\alpha) + b_{-n} \sin n(x-\alpha)) =$$

$$= b_0 + \sum_{n \in \mathbb{Z}^+} (b_n \cos nx + b_{-n} \sin nx) \cos n\alpha + (b_n \sin nx - b_{-n} \cos nx) \sin n\alpha =$$

$$= b_0 + \sum_{n \in \mathbb{Z}^+} (b_n \cos nx + b_{-n} \sin nx) \cos n\alpha + H_Z (y \mapsto b_n \cos ny + b_{-n} \sin ny)(x) \sin n\alpha.$$

Similar formula holds for $b \in \mathcal{B}^1$.)

Applying the auxiliary involution $i$ we can take complex Fourier expansion, too. This expansion can be used as follows.

9.2. The spaces $\mathcal{K}_{i,j}(\mathcal{B}, R)$ and $\Psi_{i,j}(\mathcal{B}, R)$. Let $\mathcal{K}_{i,j}(\mathcal{B}, R)$ denote the space $C_{\mathbb{Z}}(\mathcal{B}^i)$ of smooth $\mathcal{B}^i$-valued loops on $\mathbb{R}$, which are periodic for $j = 0$ and skew-periodic for $j = 1$ (in what follows: $[j]$-periodic) endowed with product operations as follows: For $b \in \mathcal{B}^j$ and $C \in \mathcal{K}_{i,j}(\mathcal{B}, R)$ we define

$$C * b = C(\beta) R^\beta b \frac{d\beta}{2\pi} \in \mathcal{B}^i.$$

So, $C$ represents a map $\mathcal{B}^j \to \mathcal{B}^i$. For $C_1 \in \mathcal{K}_{i,j}(\mathcal{B}, R)$ and $C_2 \in \mathcal{K}_{i,k}(\mathcal{B}, R)$ we set

$$C_1 * C_2(\alpha) = \int_\beta C_1(\beta) R^\beta C_2(\alpha - \beta) \frac{d\beta}{2\pi}.$$
A natural rotation action is given by
\[(R^\gamma C)(\alpha) = R^\gamma (C(\alpha)).\]

In this situation \(C\) allows a Fourier expansion
\[R^\beta C(\alpha) = \sum_{n,m \in \mathbb{Z}} c_{n,m} e^{i(-n\beta+m(\beta+\alpha))}, \quad c_{n,m} \in \mathfrak{B}^j.\]

Then \(C^*\) can be represented by the matrix \(\sum c_{n,m} e^{n,m}\). This example shows that the usual matrix computations can be transferred to these cases where “the variables are not explicit”.

If, instead of smooth loops, we use \(\mathfrak{B}^j\)-valued distributions then one can see that the multiplication operation above still exist (it is only a matter of partial integration). In that way we obtain the much larger algebra \(\Psi_{i,j}(\mathfrak{B}, R)\). What we really need is \(\tilde{K}_{i,j}(\mathfrak{B}, R)\) which is the extension of \(K_{i,j}(\mathfrak{B}, R)\) by the identity and Hilbert transform elements \(1_{Z_i} = \delta_{Z_i}^0 (\alpha)\) and \(H_{Z_i} = h_{Z_i}^0 (\alpha)\). In case of the Hilbert transforms it is particularly convenient to use matrices after we have established their properties using Lemma 8.5 or in some other ways.

9.3. Cocycles. We say that \(C \in C_i(\mathfrak{B}^0)\) is a cocycle if the identities
\[C(0) = 1 \quad \text{and} \quad C(\alpha) = C(\beta)R^\beta C(\alpha - \beta)\]
hold. We denote the set of such cocycles by \(CC_i(\mathfrak{B}^0)\).

To any invertible element \(a \in \mathfrak{B}^j\) we can associate the trivial cocycle
\[\tilde{a}(\alpha) = aR^\alpha a^{-1},\]
which is in \(CC_j(\mathfrak{B}^0)\). Not all cocycles are trivializable. For example, there may be cocycles in \(CC_1(\mathfrak{B}^0)\) even if \(\mathfrak{B}^1 = 0\). In fact, as skew-periodic cocycles are \(\mathfrak{B}^0\)-valued loops we see that skew-periodic cocycles have, a priori, nothing to do with skew-periodic loops. A second remark is that in the case of \(C\) coefficients, in general, one cannot construct a skew-periodic loop from a periodic one naturally, but one can construct a skew-periodic cocycle from a periodic one naturally.

A skew-periodic cocycle is linear, if it is of shape
\[C(\beta) = \cos \frac{\beta}{2} + G \sin \frac{\beta}{2},\]
where \(G\) is a skew-involution and \(G\) is \(R\)-invariant (i.e. \(R^\alpha G = G\) for all \(\alpha\)).

9.4. “Matrix” extension. Let \(K_{i,j}(\mathfrak{B}^k)\) contain those \(\mathfrak{B}^k\)-valued smooth functions
\[(\eta, \beta) \mapsto C(\eta, \beta)\]
which are \([i+j]\)-periodic in \(\eta\) and \([j]\)-periodic in \(\beta\), and which multiply as
\[C_1 \ast C_2 (\eta, \beta) = \int_\gamma C_1 (\eta, \gamma) C_2 (\eta + \gamma, \beta - \gamma) \frac{d\gamma}{2\pi} .\]

Then \(\Psi_{i,j}(\mathfrak{B}^k)\) can be defined similarly, these are distributions in the variable \(\beta\) and smooth in the variable \(\eta\); etc. Passing to the variable \(-\eta\) instead of \(\eta\) we see that these algebras are indeed isomorphic to matrix algebras.

It may not be obvious but these spaces also behave as tensor products: One can construct maps
\[\mathfrak{B}^{i+j+k} \otimes \Psi_{i,j}(\mathfrak{B}, R) \xrightarrow{\psi} \Psi_{i,j}(\mathfrak{B}^k)\]
where
\[ e(b \otimes A)(\eta, \beta) = b R^\eta(A(\beta)). \]

One can see that via this construction both \( \mathfrak{B}^* \) and \( \Psi_{i,j}(\mathfrak{B}) \) embed through taking tensor products with 1. (We remark, however, that \( \mathcal{K}_{1,0}(\mathfrak{B}^0) \neq 0 \) is possible even if \( \mathfrak{B}^1 \neq 0 \). Hence the image of the map above not necessarily dense.)

There are two kinds of rotations to be considered here: One is the external rotation
\[ (R^\mu C)(\eta, \beta) = R^\mu(C(\eta, \beta)), \]
and the other one is the internal rotation
\[ (\bar{R}^\mu C)(\eta, \beta) = R^\mu(C(\eta - \mu, \beta)). \]

They are related to the tensor product construction by the commutative diagrams
\[
\begin{array}{ccc}
\mathfrak{B}^{i+j+k} \otimes \Psi_{i,j}(\mathfrak{B}, R) & \longrightarrow & \Psi_{i,j}(\mathfrak{B}^k) \\
R^\gamma \otimes R^\gamma & \longrightarrow & R^\gamma \\
\mathfrak{B}^{i+j+k} \otimes \Psi_{i,j}(\mathfrak{B}, R) & \longrightarrow & \Psi_{i,j}(\mathfrak{B}^k)
\end{array}
\]

When we are interested in finding cocycles in that algebra, then the cocycle variable will be \( \alpha \) and a candidate cocycle will be given by the function \( C(\alpha, \eta, \beta) \); dependence on additional factors will be indicted in variables in front of \( \alpha \). Then the \( \Psi_{i,j}(\mathfrak{B}^0) \)-valued linear cocycles associated to the internal rotations are of shape
\[ \cos \frac{\alpha}{2} + \sin \frac{\alpha}{2} R^\gamma G(\beta), \]
where \( G(\beta) \in \Psi_{i,i}(\mathfrak{B}, R) \) is an involution.

We can formulate our general theorem: It is related to the special kernel in \( \mathfrak{g} \mathfrak{s} \mathfrak{n} \) by taking the cocycle \( c(x_1, \lambda) = \tilde{a}(x_1, x_1 - \lambda), \alpha = x_1 - x_2, \eta = x_1 - y_1, \beta = y_1 - y_2 \) and suppressing the variable \( x_1 \).

9.5. Theorem. We can define \( K_\mathfrak{g}(c, \theta, \alpha, \eta, \beta) \) such that
i.) It is a smooth perturbation of
\[ \cos \frac{\alpha}{2} \delta_2(\beta) + \sin \frac{\alpha}{2} h_2(\beta). \]
ii.) It has the cocycle property with respect to internal rotations:
\[ K(c, \theta, 0, \eta, \beta) = \delta_2(\beta), \]
\[ K_\mathfrak{g}(c, \theta, \alpha, \eta, \beta) = \int_\gamma K(c, \theta, \sigma, \eta, \gamma) R^\sigma K(c, \theta, \alpha - \sigma, \eta - \sigma + \gamma, \beta - \gamma) \frac{d\gamma}{2\pi}, \]
iii.) Value at \( \theta = 0 \):
\[ K(c, 0, \alpha, \eta, \beta) = \cos \frac{\alpha}{2} \delta_2(\beta) + \sin \frac{\alpha}{2} h_2(\beta) R^\eta c(\beta), \]
iv.) Value at \( \theta = \frac{\pi}{2} \):
\[ K(a, \frac{\pi}{2}, \alpha, \eta, \beta) = \cos \frac{\alpha}{2} \left( \delta_2(\beta) - 1 \right) + \sin \frac{\alpha}{2} h_2(\beta) + c(\alpha). \]
v.) If \( G \) is a skew-involution
\[ K(\cos \frac{\lambda}{2} + G \sin \frac{\lambda}{2}, \theta, \alpha, \eta, \beta) = \cos \frac{\alpha}{2} \delta_2(\beta) + \sin \frac{\alpha}{2} \left( h_2(\beta) + G \right). \]
vi.) It is equivariant with respect to external rotations:
\[ K_Z(R^T a, \theta, \alpha, \eta, \beta) = R^T K_Z(a, \theta, \alpha, \eta, \beta). \]

Proof. It is analogous to the matrix proof of loop linearization case but using the quadruple Fourier expansion in \( R^T, \alpha, \eta, \beta \).

9.6. In the abstract language by that we have defined a map
\[ K_Z : C^1(\mathfrak{g}) \times \mathbb{R} \to C^1(\mathfrak{g}, 0) \]
with nice properties, it linearizes (from \( \theta = \frac{\pi}{2} \) to \( \theta = 0 \)) the trivial extension of the cocycle \( c \) to the linear cocycle associated to the Bott skew-involutio
\[ J_Z(c) = \left( \beta \mapsto h^0_{\mathfrak{g}}(\beta) c(\beta) \right) \in \tilde{K}_0, 0(\mathfrak{g}, R). \]

9.7. Let us take a look to concrete example where the above theorem have meaning:
Suppose that \( P \to U \) is an smooth, oriented \( S^2 \)-bundle where the spheres come with a standard metric. Assume that \( V \) is a smooth \( K(P) \)-representative which means for us that a sufficiently nice splitting of \( C_Z = \bigoplus_{k \in \mathbb{Z}} C e_k \) is given over every point of \( P \). Assume that \( s : X \to P \) is a section, and \( V|_s(X) \) is trivial. Let \( W \) be the vertical, \( S^2 \)-tangent bundle of \( P \) over \( s(X) \).

Then on the unit circle \( S^1 W_s(u) \) of a fiber \( W_s(u) \) we can canonically construct a \( \text{GL}^\infty(C_Z) \)-valued clutching cocycle for \( V \) through exponentiation and parallel transport. Due to the orientedness of the \( S^2 \)-bundle there is a canonical \( S^1 \)-action on \( S^1 W \), hence it yields a periodic \( S^1 \)-cocycle. Due to the presence of complex numbers we can canonically construct a skew-periodic cocycle. Then, through linearization, this is homotopic to a nice \( \text{GL}^\infty(C_Z \otimes C(S^1 W_s(u))) \)-valued cocycle. This is determined by a sufficiently nice skew-involution, or rather splitting of in \( C_Z \otimes C_0(S^1 W_s(u)) \). Due to the presence of many space-dimensions we see that these are homotopic to splittings where the nontivial part of the splitting happens in \( C_Z \otimes C_1 s(u) \). We can realize that this yields the Thom isomorphism \( K(P, X) \to K(X) \) in this limited setting. In fact, this was a very special situation, we applied the linearization theorem to principal \( S^1 \)-bundles instead of general \( S^1 \)-actions, etc.

9.8. If one wants to linearize trivial cocycles (trivially) then one can define \( N(a, \theta, \eta, \beta) \) similarly.

References

[1] Pressley, A.; Segal, G.: Loop groups. Oxford Mathematical Monographs. Oxford University Press, New York, 1986.
[2] Gy. Lakos: Self-stabilization in certain infinite-dimensional matrix algebras. arXiv:math.KT/0506059