A MULTIWAVELENGTH CHARACTERIZATION OF PROTO-BROWN-DWARF CANDIDATES IN SERPENS

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Abstract

We present results from a deep submillimeter survey in the Serpens Main and Serpens/G3–G6 clusters, conducted with the Submillimetre Common-User Bolometer Array (SCUBA-2) at the James Clerk Maxwell Telescope. We have combined near- and mid-infrared spectroscopy, \textit{Herschel} PACS far-infrared photometry, submillimeter continuum, and molecular gas line observations, with the aim of conducting a detailed multiwavelength characterization of “proto-brown-dwarf” (proto-BD) candidates in Serpens. We have performed continuum and line radiative transfer modeling and have considered various classification schemes to understand the structure and the evolutionary stage of the system. We have identified four proto-BD candidates, of which the lowest-luminosity source has an $L_{\text{bol}} \sim 0.05 L_{\odot}$. Two of these candidates show characteristics consistent with Stage 0/I systems, while the other two are Stage I/I-Class Flat systems with tenuous envelopes. Our work has also revealed a $\sim20\%$ fraction of misidentified Class 0/I/Flat sources that show characteristics consistent with Class II edge-on disk systems. We have set constraints on the mass of the central object using the measured bolometric luminosities and numerical simulations of stellar evolution. Considering the available gas+dust mass reservoir and the current mass of the central source, three of these candidates are likely to evolve into BDs.

Key words: brown dwarfs – stars: formation – stars: low-mass – stars: protostars

1. INTRODUCTION

The “Serpens Main” or the “Serpens Core” cluster has been the subject of several studies at various wavelength regimes to understand the influence of the cluster environment in the process of star and planet formation. This young embedded cluster is concentrated into the northwestern (NW) and southeastern (SE) clumps or subclusters (Casali et al. 1993; Davis et al. 1999). About 45\degree south of the Serpens Main lies another active star-forming region, “Serpens/G3–G6” (Serp/G3–G6), which is a complex around four T Tauri stars named G3 to G6 in a 30′ field (e.g., Djuvpik et al. 2006; Harvey et al. 2006).

Several infrared surveys conducted in these clusters have identified more than 200 young stellar objects (YSOs) in various evolutionary stages, from prestellar cores to Class 0 to the more evolved Class II young disk sources (e.g., Evans et al. 2003, 2009; Harvey et al. 2006, 2007a, 2007b). A high population of Class 0/I systems has been found in the NW and SE subclusters, while the more evolved Class II/I/III sources are dispersed over the larger region (e.g., Kaas et al. 2004; Harvey et al. 2006, 2007a, 2007b; Winston et al. 2007). Far-infrared \textit{Herschel} observations have revealed new protostars (Bontemps et al. 2010; Könyves et al. 2010). In the submillimeter/millimeter regime, Enoch et al. (2007, 2009) identified 12 continuum sources in CSO Bolocam observations at 1 mm. A more extensive submillimeter catalog in Serpens Main has been constructed using the SCUBA bolometer array camera on the James Clerk Maxwell Telescope (JCMT), named the SCUBA Legacy Catalog (SCL; Di Francesco et al. 2008). Four prestellar cores were later identified in the SLC (Sadavoy et al. 2010). Recent observations with Combined Array for Research in Millimeter-wave Astronomy (CARMA) have identified a network of filament structures that may fragment into protostellar cores (Lee et al. 2014).

We have conducted a deep submillimeter continuum survey of the Serpens Main and Serpens/G3–G6 clusters, using the Submillimetre Common-User Bolometer Array (SCUBA-2; Holland et al. 2013) at the JCMT telescope. We have combined near-, mid-, and far-infrared observations with submillimeter continuum and molecular gas line observations, with the aim of conducting a detailed multiwavelength characterization of “proto-brown-dwarf” (proto-BD) candidates in Serpens. Previous SCUBA Legacy observations conducted in these clusters reached a (median) rms of $\sim70$ mJy/beam at 850 \textmu m (Di Francesco et al. 2008). Our survey reaches an rms that is more than $\sim10$ times deeper than the SLC, thus enabling the detection of YSOs at the very-low-mass (VLM) end that remained undetected in previous submillimeter observations. Our multiwavelength characterization has also led to the identification of YSOs that may have previously been “misidentified” as early-stage Class 0/I embedded sources, based on their near- to mid-infrared spectral index, and show characteristics that are more consistent with Class II objects with edge-on disks.

The distance to Serpens is still a matter of debate. Several studies have adopted a distance of 260 ± 10 pc from Straizys et al. (2003), based on the spectral and luminosity class of the observed stars. The 260 pc estimate is probably correct to the front of the cloud or to the front edge of the Aquila rift, but the depth of the cluster behind it is unknown. A common distance of 260 pc has been suggested for the Aquila rift and Serpens Main cluster, based on the argument that these are parts of the same star-forming region (Bontemps et al. 2010). A farther distance of 415 pc has been obtained by Dzib et al. (2010) using the Very Long Baseline Array (VLBA) parallax.
measurement of the young AeBe star EC95 in the SE subcluster of Serpens Main, Winston et al. (2010) obtained a distance of 360 pc for the Serpens core based on the X-ray luminosity function of the cluster. Their revised estimate based on new X-ray data is 260 pc for Serpens South and closer to 310 pc for Serpens North (E. Winston 2016, private communication). Another test of the distance comes from placing the spectroscopically confirmed members on the HR diagrams, and a distance of 415 pc does not correlate with the age of the cluster (E. Winston et al. 2016, in preparation). Our sample selection criteria and some of the analysis are based on the luminosity estimates from Evans et al. (2009), who had adopted a distance of 260 ± 10 pc for Serpens. We have therefore opted for a 260 pc distance as it is more suitable for this work, rather than the recent estimate of 415 pc.

The sample and various observations are described in Section 2. Results from the continuum and line radiative transfer modeling, along with a description of the bona fide detections, and the misidentified objects are presented in Section 3. A discussion on the classification of the YSOs using different schemes, estimates on the total mass of the system, and the likelihood of the proto-BD candidates evolving into a substellar object is presented in Section 4.

2. SAMPLE, OBSERVATIONS, AND DATA REDUCTION

2.1. Sample

The most detailed hunt for YSOs in the Serpens Main and Serp/G3–G6 clusters was conducted under the Spitzer “Cores to Disks” (c2d) legacy project (Evans et al. 2003). We searched the c2d catalog for all early-stage (Class 0/1/Flat) YSOs with bolometric luminosity \( L_\text{bol} \lesssim 0.3 L_\odot \). We selected this limit since it corresponds to the typical boundary between VLM and low-mass objects (e.g., Chabrier & Baraffe 1997, 2000). Considering the mass–luminosity relations based on the nonaccretion models by Baraffe et al. (2003), an \( L_\text{bol} \sim 0.3 L_\odot \) would correspond to a stellar mass of \( \sim 0.2–0.3 M_\odot \), assuming an age of \( \sim 1 \) Myr for the cluster. Based on the same models, substellar mass objects (\( M_s \approx 0.08 M_\odot \)) or BDs at this age would have \( L_\text{bol} \approx 0.05 L_\odot \). Our target list consists of 28 VLM/BD objects, 13 of which are classified as Class 0/1 and 15 as Class Flat sources, as listed in the Spitzer c2d catalog.

It is important to note a possible uncertainty in converting \( L_\text{bol} \) into stellar mass using nonaccreting substellar evolutionary models, as in Baraffe et al. (2003), which do not include the contribution to the luminosity from ongoing accretion. The bolometric luminosity in such a case can only provide an upper limit to the stellar luminosity or mass of the source. We have taken into account the accretion effects and set constraints on the mass of the central object using new, more accurate accreting models, as presented in Section 4.2.2.

2.2. Submillimeter Continuum Observations

We obtained submillimeter continuum observations using the JCMT SCUBA-2 (Holland et al. 2013) instrument. SCUBA-2 provides dual-wavelength observations at 450 and 850 \( \mu m \). The default map pixels are 2″ and 4″, and the half-power beam width is 7″/5 and 14″/5 in the 450 and 850 \( \mu m \) bands, respectively. The observations were obtained in the months of 2014 April–June (PID: M14AU05) in Grade 2 weather (225 GHz opacity of 0.06). We used the CV Daisy observing mode and also applied a matched-beam filter that utilizes the full flux in the beam. Using this setup, we obtained a 1σ rms of \( \sim 3 \) mJy/beam at 850 \( \mu m \). The reduction and calibration of the SCUBA-2 data were conducted using the Sub-Millimetre User Reduction Facility (SMURF), which is available under the Starlink package. The final science maps were produced using the various steps and the default configurations described in the map-making process Dynamic Iterative Map-Maker (Chapin et al. 2013). We used the KAPPA and PICARD software for the postprocessing of the science maps. To improve the point-source detectability, the PICARD recipe SCUBA2_MATCHED_FILTER was applied to the flux-calibrated science maps. This recipe fits a single Gaussian point-spread function (PSF), centered over every pixel in the map, and applies a background-suppression filter to remove any residual large-scale noise. The half-power beam width in the respective bands was used as the full width at half maximum of the Gaussian fit. The images are convolved with the modified PSFs. This technique is recommended for faint source extraction and has been utilized before in, for example, creating the SLC (Di Francesco et al. 2008). The matched filtering technique is valuable in extracting the optimum flux from such low-luminosity objects and helps to minimize the contamination from the surrounding cloud material. The full scale of the 850 \( \mu m \) maps is shown in Figure 1.

2.3. Near-infrared Spectroscopy

We obtained near-infrared long-slit spectroscopy with the LIRIS spectrograph on the William Herschel Telescope in 2015 May. We used the moderate-resolution (\( R \sim 2500 \)) \( R_h \), \( R_h \), and \( R_k \) grisms with a slit width of 1″ and a two-nod-points setup. The seeing varied between 0″8 and 1″0. The beginning parts of all observing nights were lost because of cloudy weather conditions. Due to the time constraint, we were only able to obtain \( K \)-band spectra for five sources. The spectra were reduced using the IRAF IAC irisdr task package. This package includes the standard routines for bias subtraction, flat fielding, wavelength, and flux calibration. The wavelength scale was corrected to the heliocentric standard of rest. We estimate a signal-to-noise ratio (S/N) of \( \sim 5–10 \) in the whole spectra. The near-infrared spectra are shown in Figure 2.

2.4. Mid-infrared Spectroscopy

We obtained mid-infrared low-resolution (\( R \sim 600 \)) \( N \)-band spectroscopy with the CanariCam spectrograph on the Gran Telescopio Canarias in 2015 June. The observations were conducted under conditions with PWV ≤ 8 mm. We observed using the chop-nod mode and a slit width of 0″5. We obtained a 1σ point-source sensitivity of 55 mJy at 10 \( \mu m \). For data reduction, we used the RedCan pipeline developed by Gonzalez-Martin et al. (2013). The final products from this pipeline are flux-calibrated spectra combined from multiple exposures of each science target. We estimate an S/N > 10 for the spectra; the S/N is lower at the very base of the spectrum near \( \sim 9.5 \mu m \). For three sources, low-resolution Spitzer/IRS spectra were available in the archives, and these were extracted and calibrated using the Spitzer IRS Custom Extraction software. Further details on the processing of the Spitzer/IRS spectra are provided in Riaz (2009). The mid-infrared spectra are shown in Figure 3.
2.5. Archival Data

2.5.1. Herschel PACS Observations

We have analyzed the archival Herschel PACS 70, 100, and 160 \( \mu m \) scans of the Serpens Main and Serp/G3–G6 clusters. These observations were taken from the programs ObsIDs 134229080 and 1342206676. We analyzed the final pipeline-processed “Level 2.5” data and used the rectangular-sky photometry task provided in the Herschel Interactive Processing Environment (HIPE). The rectangular photometry task applies aperture photometry to both a circular target aperture and a rectangular sky aperture. The sky intensity was estimated using the median-sky estimation algorithm. Since the sky intensity could vary with location, we selected four different background regions around the target and measured the sky-subtracted source flux in each sky region. The final source flux and error are the mean and standard deviation of these four flux measurements.

2.5.2. Molecular Line Observations

We searched for molecular line observations in various archives, and we were able to retrieve observations of the HCO\(^+\) (3-2) transition obtained with the CSO heterodyne receiver (XFFTS spectrometer). The final processed CLASS

![Figure 1. SCUBA-2 850 \( \mu m \) maps in THE Serpens Main (top) and Serp/G3–G6 (bottom) clusters. The color scale at the bottom shows the intensity in units of Jy beam\(^{-1}\). North is up, and east is to the left.](image)
files for the HCO\(^+\) data were kindly sent by N. Evans and A. Heiderman, and some of the data were obtained from the CSO archives. The data-reduction process of the HCO\(^+\) data is discussed in Heiderman & Evans (2015). There were no C\(^{18}\)O line observations available in the archives for the VLM BD sources detected in the SCUBA-2 maps. The molecular line spectra for the individual sources are discussed in Section 3.4. In addition, there are six proto-BD candidates that are undetected (S/N < 2) in the SCUBA-2 maps but show emission in the HCO\(^+\) (3–2) molecular line (Section 3.5). The discussion on sources with marginal submillimeter detection (S/N ∼ 2) and sources that are unresolved or confused in the SCUBA-2 and PACS maps due to close proximity to a more luminous protostar is presented in the Appendix.

3. RESULTS

We have detected seven Class 0/I and eight Class Flat very-low-luminosity YSOs in our SCUBA-2 submillimeter observations. A detailed description of the YSOs that have a bona fide detection in the 850 μm band, with an S/N ≥ 5, is provided in Section 3.4. Among the bona fide detections, there are four “misidentified” cases for which the evolutionary stage based on the physical characteristics appears inconsistent with their original spectral energy distribution (SED) class provided in the c2d catalog (Section 3.4.1). In addition, there are six proto-BD candidates that are undetected (S/N < 2) in the SCUBA-2 maps but show emission in the HCO\(^+\) (3–2) molecular line (Section 3.5). The discussion on sources with marginal submillimeter detection (S/N ∼ 2) and sources that are unresolved or confused in the SCUBA-2 and PACS maps due to close proximity to a more luminous protostar is presented in the Appendix.

3.1. Near- and Mid-infrared Spectra

The near- and mid-infrared spectra provide an independent diagnostic of the evolutionary stage of the system. The mid-infrared spectra for five sources observed show a deep absorption feature (Figure 3), with a peak at ~9.8 μm indicative of amorphous silicate dust and a shape similar to the ones typically observed in Class I protostellar sources (e.g., Kessler-Silacci et al. 2005), while one source, J182902.84+003009.6, shows a combined absorption+emission feature indicative of an edge-on disk (Section 3.4.1). Another important youth indicator is near-infrared spectra. The quality of the near-infrared spectra is not high enough to conduct a detailed analysis of the accretion and outflow properties in these sources. The main purpose of obtaining near-infrared spectra was to determine the youth of the sources, in particular to confirm that the spectra do not show profiles similar to active galactic nuclei (AGNs) or extragalactic sources, as was found to be the case for two candidates in our σ Orionis study (Riaz et al. 2015). None of the YSOs studied in this work with near-infrared spectra show profiles similar to AGNs or extragalactic sources. The near-infrared spectra (Figure 2) show strong veiling, with \( r_K > 2 \) for all sources observed. Large optical or near-infrared veiling is commonly observed in low-mass protostars and has been suggested as one of the diagnostics to distinguish between Class I and Class II sources (e.g., Connelley & Greene 2010). However, for misidentified objects such as J182956.67+011239.2 (Section 3.4.1), the presence of an edge-on disk can also result in high extinction. Large veiling is still a useful YSO indicator showing that the spectra are dominated by excess continuum emission from the circumstellar material. A discussion on the near- and mid-infrared spectra for the individual sources is presented in Section 3.4.

3.2. Radiative Transfer Modeling of the SEDs

Radiative transfer modeling of the SEDs can map the physical structure of the system and is important in obtaining estimates on the masses and sizes of the envelope and disk components, the envelope mass infall rate, and the inclination angle of the system. It can also provide better constraints on the contribution to the total luminosity from the envelope and disk components. We have built the SEDs for the bona fide detections by compiling photometric data in the near infrared from the UKIDSS and the 2MASS surveys and in the mid-infrared from the Spitzer c2d catalog, the PACS far-infrared photometry, and the SCUBA-2 submillimeter photometry. For six sources, the mid-infrared spectra have also been included in the fits. For three sources, the 1.1 mm photometry from CSO...
Bolocam observations (Enoch et al. 2009) is available. However, the 1.1 mm flux density is higher than the 850 μm photometry for all three sources, so we have opted not to use the 1.1 mm point in fitting the SEDs. Any model that provides a good fit to the 1.1 mm point misses the far-infrared 70–160 μm and the submillimeter 450 and 850 μm points and is a poor fit to the overall SED. The possible reasons for the discrepancy between the 850 μm and 1.1 mm photometry are discussed in Appendix A.3.

The SED modeling was conducted using the two-dimensional radiative transfer code by Whitney et al. (2003). The main ingredients of the model are a rotationally flattened infalling envelope, bipolar cavities, and a flared accretion disk in hydrostatic equilibrium. For the circumstellar envelope, the angle-averaged density distribution varies roughly as $\rho \propto r^{-1/2}$ for $r \ll R_\text{c}$, and $\rho \propto r^{-3/2}$ for $r \gg R_\text{c}$. Here, $R_\text{c}$ is the centrifugal radius and is set equal to the disk outer radius. This includes only the current infalling part of the envelope. Beyond $r_{\text{infall}}$, the density decreases as $r^{-2}$, and most of the mass is out in that region for the early (Stage 0) phase. The disk density is proportional to $\varpi^{-\alpha}$, where $\varpi$ is the radial coordinate in the disk midplane and $\alpha$ is the radial density exponent. The disk scale height increases with radius, $h = h_0(\varpi/R_*)^{\beta}$, where $h_0$ is the scale height at $R_*$ and $\beta$ is the flaring power. The disk extends from the dust sublimation radius, $R_{\text{sub}} = R_*(T_{\text{sub}}/T_*)^{-2/3}$, to the outer disk radius, $R_{\text{disk, max}}$. Here, $T_{\text{sub}}$ is the dust sublimation temperature, which was set to 1600 K. The bipolar cavities in the models extend

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**Figure 3.** Mid-infrared spectra. The vertical lines mark the amorphous olivine and crystalline forsterite peaks at 9.8 μm and 11.3 μm, respectively.
from the center of the protostar to the envelope outer radius ($R_{\text{max,env}}$). We adopted the curved-shape cavity, which has the structure of $z = ax^2$, where $x = (\chi^2 + \chi^2)^{1/2}$ and the constant $a$ is determined by a relation between the envelope radius and the cavity opening angle (Whitney et al. 2003).

Table 3 lists the estimates for the various envelope and disk model parameters, based on the best model fit (lowest $\chi^2$ value) to the observed SED. The uncertainties for each parameter represent the range in values obtained from the degeneracies in the top five model fits. The parameters that notably affect the model SEDs are the mass infall rate ($M_{\text{env}}$), the inclination angle ($\theta_{\text{in}}$), and the centrifugal radius. An increase in $M_{\text{env}}$ corresponds to a denser envelope, which implies that a lesser number of photons can escape through the cavity regions, thus corresponding to a denser envelope, which implies that a lesser number of photons that can escape through the cavity regions, thus producing lower near-infrared fluxes. The best model fits to the observed SEDs for the Class 0/I targets were obtained for $M_{\text{env}} \sim 10^{-5} - 10^{-6} M_\odot \text{yr}^{-1}$ (Table 3). Higher values for $M_{\text{env}}$ result in a model with excess emission in the submillimeter regime and do not fit the 850 μm point. We can therefore constrain the $M_{\text{env}}$ parameter using both the submillimeter and the near-infrared points.

The amount of flaring in the system, particularly around ~100 μm, can be controlled with the centrifugal radius. For smaller values of $R_c$, the model SEDs show more flux at wavelengths >24 μm and lower fluxes below ~8 μm. Decreasing $R_c$ implies that infalling material piles up closer to the central protostar. This results in fewer photons that can escape through the cavity walls and the disk upper layers and shifts the far-infrared peak in the model SED toward shorter wavelengths. Likewise, the effects of the inclination angle on the model SED are also mainly seen in the far infrared. A high-inclination system implies that there is more absorbing material in our line of sight. Thus with decreasing inclination angle, there is a decline in the far-infrared fluxes near 100 μm and an increase in the near- and mid-infrared fluxes. The effects of increasing the outer radius of the envelope ($R_{\text{env, max}}$) are mainly seen in the model SED at the longest wavelengths. This parameter was constrained by the far-infrared/submillimeter points, and a value of ~1500–2000 au provides a good fit. The depth, shape, and width of the 10 μm silicate absorption feature provide additional constraints on the envelope and cavity parameters in the model fits. Increasing the envelope density results in a deeper silicate absorption profile, whereas increasing the cavity opening angle results in a shallower silicate feature. Among the disk parameters, the outer disk radius ($R_{\text{disk, max}}$) was set to the same value as the centrifugal radius, while the inner disk radius is a few stellar radii (~3–5 $R_{\text{sub}}$). Results from the SED modeling for the individual sources are discussed in Section 3.4.

### 3.3. Radiative Transfer Modeling of the Molecular Lines

The HCO$^+$ (3-2) transition can probe the dense gas located in the inner regions of protostellar envelopes and is known to be an important indicator of a Class 0/I source (e.g., van Kempen et al. 2009). The interpretation of HCO$^+$ data requires a simple model to verify the dynamics of the system. We adopt the “two-layer” model of Di Francesco et al. (2001), based on Myers et al. (1996). The two layers are assumed to be a uniform slab layer characterized by a single temperature $T_f$ and $T_r$ for the front and rear layers, respectively. The difference from the earlier Myers et al. (1996) model is the addition of a continuum layer $T_c$ in between. The front layer is a representative of the redshifted layer, which has H$_2$ number densities between 10$^4$ and 10$^6$ cm$^{-3}$. An estimation for the excitation temperature with an HCO$^+$ column of 10$^{12}$–10$^{13}$ cm$^{-2}$ yields 7–9 K, based on RADEX models. This is subthermally excited gas assuming the emitting region is between 15 and 20 K. On the other hand, the rear layer is always thermalized at between 15 and 20 K depending on the physical structure. These values were determined from the typical temperatures at >500 au in the system, considering the best-fit SED models. Here, we fixed the front layer to 9 K and a rear layer to 15 K. The free parameters are infall velocities $\upsilon_{\text{infl}}$, optical depth $\tau_0$, and the filling factor $\Phi$. The models also have a broadening factor for the Gaussian line profile. The exact values may change depending on the adopted excitation temperatures. However, the goal of this modeling is to characterize the infalling velocities, which showed a weaker dependence on the adopted excitation temperatures. The systemic cloud velocity in Serpens is taken to be ~8 km s$^{-1}$ (e.g., Kirk et al. 2013). The model profiles are further discussed in Sections 3.4 and 4.1.

### 3.4. Bona Fide Detections

$\text{SSTc2d J182902.12+003120.7}$ ($\text{J182902.12}$): This is a Class Flat ($\alpha_{\text{fl}} = 0.27$) source with an $L_{\text{bol}} = 0.05 L_\odot$ (Table 1). The SCUBA-2 maps show a clearly extended object in the 450 μm

### Table 1

Object Positions and Properties of the Bona Fide Detections

| Object Name | R.A. [J2000] | Decl. [J2000] | $L_{\text{bol}}$ | $\alpha_{\text{fl}}$ | $T_{\text{bol}}$ | HCO$^+$ Detection | Class$^c$ | Revisedd |
|-------------|--------------|--------------|-----------------|--------------------|-----------------|-----------------|----------|--------|
| J182902.12+003120.7 | 18:29:02.1 | +00:31:20.7 | 0.05 | 0.24 | 85 | Y | Flat, I, 0/1 | Flat, I/T |
| J182855.78+002944.8 | 18:28:55.7 | +00:29:44.8 | 0.18 | 1.89 | 54 | Y | 0, 0, 0/1 | 0, 0 |
| J182949.57+011706.0 | 18:29:49.5 | +01:17:06.0 | 0.25 | 0.66 | 480 | NA | I, L, – | I, T |
| J182841.87+000321.3 | 18:28:41.8 | +00:03:21.3 | 0.16 | 0.37 | 410 | N | 1, Flat, II | Flat, II |
| J182956.67+011439.2 | 18:29:56.6 | +01:14:39.2 | 0.06 | –0.11 | 690 | Y | Flat, Flat, 0/1 | II, II |
| J182902.84+003009.6 | 18:29:02.8 | +00:30:09.6 | 0.17 | –0.14 | 460 | NA | Flat, Flat, – | II, II |
| J182955.69+011431.6 | 18:29:55.6 | +01:14:31.6 | 0.2 | –0.23 | 810 | NA | Flat, Flat, – | II, II |

#### Notes:

$^a$ The estimates for $L_{\text{bol}}$, $\alpha_{\text{fl}}$, and $T_{\text{bol}}$ are from Evans et al. (2009).

$^b$ Y—Yes; N—No; NA—HCO$^+$ (3-2) observations not available.

$^c$ The first and second values indicate the classification using the schemes based on $\alpha_{\text{fl}}$ and $T_{\text{bol}}$, respectively, from the criteria defined in Evans et al. (2009). The third value is the classification using the criteria based on the HCO$^+$ line strength from Heiderman & Evans (2015).

$^d$ The first value indicates the revised classification based on SED modeling results (Section 3.2); the second value is obtained using the “Stage” classification scheme.

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band but less extended emission at 850 μm (Figure 4). There is also a clear detection in the PACS bands. It has been previously detected in the dust continuum emission at 1.1 mm (Enoch et al. 2009), based on which it was categorized as an envelope source (Evans et al. 2009). J182902.12 is undetected in the near-infrared bands.

The SED for J182902.12 shows deep absorption in the mid-infrared silicate feature, followed by a steep rise in the far infrared (Figure 5). An envelope component is required in order to fit the far-infrared and submillimeter points. The best model fit indicates an envelope mass of ~8 $M_{\text{Jup}}$ and a disk mass of ~2 $M_{\text{Jup}}$, with an edge-on inclination of ~80° for the system (Table 3). The HCO$^+$ spectrum shows weak emission with a peak at the cloud velocity (Figure 6).

**Table 2**
Flux Densities of the Bona Fide Detections

| Object Name | PACS | SCUBA-2 | 1.1 mm |
|-------------|------|---------|--------|
| (SSTc2d +)  | 70 μm (mJy) | 100 μm (mJy) | 160 μm (mJy) |
| 70 μm (mJy) | 100 μm (mJy) | 160 μm (mJy) |
| 100 μm (mJy) | 100 μm (mJy) | 160 μm (mJy) |
| 160 μm (mJy) | 160 μm (mJy) | 160 μm (mJy) |

Note. $^a$ The 1.1 mm flux densities are from Enoch et al. (2009).

The best SED model fit for J182902.12 shows a massive envelope of ~0.2 $M_\odot$ and a disk component of ~7 $M_{\text{Jup}}$ at a ~30° inclination for the system (Figure 5; Table 3). The HCO$^+$ (3-2) emission shows a broad profile, with a self-absorption component at the cloud systemic velocity and a hint of a red-dominated asymmetry (Figure 6).

SSTc2d J182855.78+002944.8 (J182855.78): Among the bona fide detections, J182855.78 is the most deeply embedded Class 0 source ($\alpha_{\text{IR}}$ = 1.89), with a $L_{\text{bol}}$ = 0.18$L_\odot$ (Table 1). This is a bright extended object in the SCUBA-2 maps (Figure 4). The extended shape is more clearly seen in the 450 μm map. There is a point-source detection in the PACS 70 and 100 μm bands but a marginal detection in the 160 μm. This source is undetected in the near infrared and is too faint for mid-infrared spectroscopy.

The best model fit to the SED shows a massive envelope of ~0.2 $M_\odot$ and a disk component of ~7 $M_{\text{Jup}}$ at a ~30° inclination for the system (Figure 5; Table 3). The HCO$^+$ (3-2) emission shows a broad profile, with a self-absorption component at the cloud systemic velocity and a hint of a red-dominated asymmetry (Figure 6).

SSTc2d J182841.87-000321.3 (J182841.87): This is a Class I source ($\alpha_{\text{IR}}$ = 0.66) but is characteristic of VLM YSOs of mid-M spectral types (e.g., Connelley & Greene 2010; Gorlova et al. 2010). The K-band spectrum shows weak detection in the $H_2$ (1-0) $S$ line at ~2.05 μm (Figure 2), which is a known indicator of outflow activity in protostars (e.g., Caratti o Garatti et al. 2006), but no other notable accretion or outflow features are seen.

The SED modeling results for J182841.87 indicate two possible fits (Figure 8). The best-fit model provides a good fit to the submillimeter points as well as the full depth of the 10 μm silicate feature but misses the mid- and far-infrared points. The second model fits the mid- and far-infrared photometry but misses the submillimeter points as well as the silicate feature. This indicates that a small envelope component is required in order to fit the submillimeter points. The best fit is obtained for an envelope mass of ~13 $M_{\text{Jup}}$, a disk mass of ~6 $M_{\text{Jup}}$, and an inclination of 40° for the system. In contrast, the second fit is mainly dominated by the disk emission and is obtained for a close to edge-on disk at an inclination of ~69° and a mass of ~7 $M_{\text{Jup}}$. There is no HCO$^+$ (3-2) line detection. The characteristics for J182841.87 are consistent with either a Stage I edge-on disk source or a Stage I/II transition object with a tenuous envelope. This object has been misidentified as an embedded Class I system.

SSTc2d J182955.69+011431.6 (J182955.69): This is a Class Flat ($\alpha_{\text{IR}}$ = −0.11) source with observed $L_{\text{bol}}$ = 0.06$L_\odot$. There is a point-source detection in the SCUBA-2 maps, but no clear point source is seen in the PACS maps (Figure 7). The PACS photometry should thus be considered as the upper limits obtained at the target position. This object is undetected in the $JH$ bands and is a very faint source in the $K$ band ($K = 14.06$ mag). This YSO shows the strongest detection in the $H_2$ (1-0) $S$ line at ~2.05 μm (Figure 2), indicating outflow activity.
All possible SED model fits for this source indicate a disk-dominated object at an edge-on inclination of ~80° with a disk mass of ~12 \( M_{\text{Jup}} \) (Figure 8; Table 3). There is an indication of an extremely tenuous envelope component of ~0.4 \( M_{\text{Jup}} \) mass. The HCO\(^+\) (3-2) line shows a broad profile centered close to the cloud velocity, with a hint of self-absorption at higher velocities (Figure 6). The Stage 0/I classification for J182956.67 by Heiderman & Evans (2015), based on HCO\(^+\) detection, is inconsistent with the SED modeling results. It may be the case that there are different kinds of sources in the beam, such as a faint star superposed on a starless core, resulting in a strong detection in the HCO\(^+\) line. An \( \alpha_R \) of ~0.11 places it closer to the Class Flat/Class II boundary (~0.3). It may possess a very tenuous envelope, if any, and is more likely to be a case of a Stage II or Class II edge-on disk source.

SSC2d J182902.84+00309.6 (J182902.84): This is a Class Flat \( (\alpha_R = -0.14) \) source with observed \( L_{\text{bol}} = 0.17 L_{\odot} \). There is a point-like detection for this object in the SCUBA-2 and PACS 70 and 100 \( \mu \)m bands, though it appears (slightly) extended in the SCUBA-2 450 \( \mu \)m map (Figure 7). There is no detection in the PACS 160 \( \mu \)m band. A detection at 1.1 mm has been previously reported, based on which it was categorized as an envelope source (Enoch et al. 2009; Evans et al. 2009). The near-infrared spectrum for this source is similar to the mid-M type seen in Class I/Flat systems (e.g., Gorlova et al. 2010) (Figure 2).

J182902.84 shows an interesting 10 \( \mu \)m silicate feature, with both an emission and an absorption component (Figure 3). Such a feature has been earlier noted in a BD edge-on disk (e.g., Luhman et al. 2007). The silicate emission component has an origin in the optically thin surface layers of the disk, whereas the absorption is caused by the highly extinguished inner disk wall. The best model fit to the J182902.84 SED shows an edge-on disk at a ~75° inclination, with a mass of ~11 \( M_{\text{Jup}} \) (Figure 8). No envelope component is required, and a disk alone is adequate to fit the submillimeter points. There are no HCO\(^+\) (3-2) observations available for this source. Overall, J182902.84 appears to be a case of a misidentified object, showing characteristics that are more similar to a disk-only Stage II or Class II source.

SSC2d J182955.69+011431.6 (J182955.69): This is a Class Flat source \( (\alpha_R = -0.23) \), with an observed \( L_{\text{bol}} = 0.2 L_{\odot} \). The spectral type determined from near-infrared spectroscopy is M4.5 ± 0.5 (Gorlova et al. 2010). J182955.69 shows a bright point-like detection in the SCUBA-2 maps, the first submillimeter detection for this object (Figure 7). It is difficult to confirm a detection in the PACS map as it lies in a nebulos, confused region. Thus the PACS photometry should be considered as upper limits. The results from SED modeling for J182955.69 indicate an edge-on disk at ~75° inclination, with a disk mass of ~18 \( M_{\text{Jup}} \) (Figure 8). The best fit also indicates a very weak envelope component of ~0.1 \( M_{\text{Jup}} \). The spectral slope lies close to the Class Flat/Class II boundary. J182955.69 is thus a misidentified object and appears as a disk-dominated Stage II or Class II source.

### Table 3

| Object Name | \( M_{\text{env}} \) \((M_{\text{Jup}})\) | \( M_{\text{disk}} \) \((M_{\text{Jup}})\) | \( M_{\text{bol}} \) \((M_{\text{Jup}})\) | \( R_{\text{env,max}} \) \((\text{au})\) | \( R_{\text{d,max}} \) \((\text{au})\) | \( \theta_{\text{PA}} \) \((\text{°})\) |
|-------------|---------------------|---------------------|---------------------|--------------------|--------------------|----------------------|
| SSTc2d J182902.12+003120.7 | 8 ± 3 | 2 ± 1 | \((0.3 ± 0.1) \times 10^{-6}\) | 1510 ± 100 | 102 ± 30 | 80 ± 10 |
| J182855.78-002944.8 | 216 ± 10 | 7 ± 2 | \((2 ± 1) \times 10^{-5}\) | 1700 ± 100 | 47 ± 5 | 32 ± 10 |
| J182949.57-011706.0 | 33 ± 10 | 17 ± 8 | \((2 ± 1) \times 10^{-6}\) | 2230 ± 500 | 10 ± 5 | 40 ± 10 |
| J182841.87-003031.3 | 13 ± 6 | 6 ± 1 | \((0.4 ± 0.2) \times 10^{-6}\) | 1420 ± 400 | 46 ± 6 | 40 ± 10 |
| J182956.67+011392.9 | 0.4 ± 0.3 | 12 ± 4 | \((2 ± 1) \times 10^{-9}\) | ... | 73 ± 20 | 80 ± 10 |
| J182902.84+003009.6 | <10^{-6} | 11 ± 3 | ... | ... | 11 ± 3 | 75 ± 10 |
| J182955.69+011431.6 | 0.1 ± 0.05 | 18 ± 5 | \((0.2 ± 0.1) \times 10^{-9}\) | ... | 90 ± 20 | 75 ± 10 |

Note. \( M_{\text{env}} \) and \( M_{\text{disk}} \) are the (gas+dust) masses.

#### 3.5. Possible Proto-BD Candidates

For six sources in our sample, there is no detection \((S/N < 2\sigma)\) in the SCUBA-2 maps. None of these sources are detected in the PACS maps. The observed \( L_{\text{bol}} \) for these objects is <0.04 \( L_{\odot} \), implying that the envelope masses are too low, and the nondetection can be explained by the poor sensitivity in both the SCUBA-2 and PACS maps. We consider these objects as potential proto-BD candidates.

Among these candidates, four sources show emission in the HCO\(^+\) (3-2) line (Figure 9). The line emissions for three sources show a broad profile with self-absorption at the systemic velocity. For the source J182852.76+002846.8, the HCO\(^+\) emission shows slightly blueshifted asymmetry indicative of an infalling envelope. A weak indication of a blue component is also seen for J182959.03+011225.1, for which the main emission with a peak at ~8 km s\(^{-1}\) must be dominated by the foreground cloud material. The line appears to be a blend of multiple velocity components, possibly arising from a molecular outflow. In contrast, the source J182947.01+011626.9 shows a slightly red-dominated asymmetry, as observed in some disk-dominated objects (e.g., van Kempen et al. 2009), although such asymmetries could also appear in some embedded sources driving an outflow (e.g., D. Harsono et al. 2016, in preparation). All four of these sources have been categorized as Stage 0/1 protostars due to HCO\(^+\) line detection (Heiderman & Evans 2015). We note that while Heiderman & Evans (2015) have searched for HCO\(^+\) observations within 14° of the Spitzer c2d source position, all four of these sources with HCO\(^+\) line detection are located at a ~20° separation from a bright source in the SCUBA-2 maps. Since the nearby bright object is within the CSO 30° beam size, it may have contaminated the observed line emission. Spatially resolved observations at higher sensitivities can provide better insight into their characteristics.
Figure 4. Herschel and SCUBA-2 images for the bona fide sources. Target position is marked by a cross. From top to bottom, the images are at 70, 100, 160, 450, and 850 μm. The sources from left to right are J182902.12, J182855.78, and J182949.57. The representative color scale at the bottom shows the intensity in units of Jy beam⁻¹. The contours are given in steps of one from 1× to 5× the peak intensity. The pixel size in the PACS 70 and 100 μm bands is 3″ and 6″ in the 160 μm band. The default map pixels are 2″ and 4″ at 450 μm and 850 μm, respectively. The spatial scale is shown at the bottom right. North is up, and east is to the left.
Section 3.4.1, we have found four Class Flat and Class II sources. Except for the transition cases close to the boundary between components, respectively.

SEDs with model fits for the bona fide sources. Red, green, and blue lines indicate the individual contribution from the envelope, disk, and stellar components, respectively.

4. DISCUSSION

4.1. Classification of VLM/Substellar YSOs

The traditional method of determining the evolutionary class of a YSO is to measure the near- to mid-infrared spectral index, \( \alpha_{IR} = d \log(\lambda F_\lambda)/d \log(\lambda) \), as first defined by Lada & Wilking (1984) and Adams et al. (1987). The typical wavelength range considered to measure \( \alpha_{IR} \) is 2–24 \( \mu \)m. The now widely used classification scheme consists of five classes of YSOs: Class 0 is the earliest, deeply embedded stage with a high submillimeter-to-bolometric luminosity ratio (Andre et al. 1993); Class I has more evolved, embedded YSOs with \( \alpha_{IR} > 0.3 \); the “Flat Spectrum” sources (\(-0.3 < \alpha_{IR} < 0.3\); Greene et al. 1994) are at an intermediate stage between Class I and II and have tenuous envelopes compared to Class I objects; Class II (\(-2 < \alpha_{IR} < -0.3\)) are T Tauri stars with gas-rich circumstellar disks but no envelope material; and Class III (\(\alpha_{IR} < -2\)) has pre-main-sequence stars that may possess tenuous disk material. Another classification method is based on the bolometric temperature, \( T_{bol} \), which was first connected to the classes defined by \( \alpha_{IR} \) by Chen et al. (1995). In Table 1 (column 8), we have listed the classification as originally determined by Evans et al. (2009), using both the spectral index and \( T_{bol} \) criteria. The classes from the two schemes are similar, except for the transition cases close to the boundary between Class Flat and Class II sources.

From our analysis based on SED modeling presented in Section 3.4.1, we have found four out of seven bona fide detections to be misidentified objects, such that their characteristics do not comply with the original classification based on \( \alpha_{IR} \) or \( T_{bol} \) criteria. These include three Class Flat sources that show characteristics similar to Class II objects and a Class III system that appears to be a Class Flat source. The revised classification based on the SED modeling results are listed in Table 1 (column 9). As noted by Whitney et al. (2003), there can be a wide range in \( \alpha_{IR} \) for a given \( T_{bol} \) category, and therefore classifying a YSO only based on its observed characteristics of the spectral slope or the bolometric temperature can be erroneous, particularly for cases such as a face-on embedded YSO, which may be misidentified as Class II, or an edge-on flaring disk, which may be misidentified as a Class I object.

The “Stage” classification scheme, introduced by Whitney et al. (2003) and later modified by Robitaille et al. (2006), is based on the evolutionary stage of a YSO using its true physical characteristics, such as the disk mass and envelope accretion rate, regardless of the observed properties. In this classification scheme, Stage 0/1 objects have significant infalling envelopes, with \( M_{env}/M_{star} > 10^{-6} \) yr\(^{-1}\); Stage II sources have tenuous envelopes but gas-rich optically thick disks, with \( M_{env}/M_{star} < 10^{-6} \) yr\(^{-1}\) and \( M_{disk}/M_{star} > 10^{-6} \); and Stage III objects with \( M_{env}/M_{star} < 10^{-6} \) yr\(^{-1}\) and \( M_{disk}/M_{star} < 10^{-6} \) have no envelopes but may have tenuous disks. Some additional constraints are as follows: Stage 0 objects have \( M_{disk}/M_{env} \ll 1 \) and \( M_{circum}/M_{star} \sim 1 \), Stage I have \( 0.1 < M_{disk}/M_{env} < 2 \) and
Figure 6. HCO$^+$ (3-2) observed (left) and model (right) line profiles for the bona fide detections. The vertical line marks the systemic cloud velocity.
M_{circum} < M_{star} and Stage II have $M_{env} = 0$ and $M_{disk}/M_{star} \ll 1$. Here, $M_{circum}$ refers to the envelope+disk mass. An intermediate phase is the Stage I Transition (Stage I-T) category (e.g., van Kempen et al. 2009), with $M_{env}/M_{star} \sim (1-2) \times 10^{-6} \, \text{yr}^{-1}$ and $M_{disk}/M_{env} > 2$. These sources have tenuous envelopes compared to the early Stage I objects.

More recently, a classification scheme combining dust continuum and molecular line observations, notably the high gas density tracer HCO$^+$ molecule, has been developed to identify the truly embedded Class 0/I sources and separate them from Class II disk sources (e.g., van Kempen et al. 2009). In addition to a submillimeter dust continuum detection peaking on the source, this criterion is based on the strength of the HCO$^+$ emission; in particular, the integrated intensity of the HCO$^+$ (4-3) line being >0.4 K km s$^{-1}$ is considered to be a good metric for a Class 0/I or Stage 0/I sources, while the
HCO\(^+\) line is either undetected or detected at a very weak level \((T_{\text{mb}} < 0.1 \text{ K})\) in disk-dominated Class II/Stage II sources. Another notable feature is an axisymmetric profile observed in the HCO\(^+\) line, which may be indicative of an infalling envelope. Typically, a blue-dominated infall asymmetry has been seen in the HCO\(^+\) profile of some low-mass protostars, while a red-dominated axisymmetry is usually associated with disk-only objects or the presence of an outflow \((\text{e.g., Thi et al. 2004; Evans et al. 1999; van Kempen et al. 2009})\). However, this is not always the case, since red-dominated asymmetry has been observed in embedded sources \((\text{D. Harsono et al. 2016, in preparation})\), an interpretation of which could be an expanding envelope that is due to an outflow.

Among the bona fide detections in our present work, J182902.12, J182855.78, and J182956.67 have previously been categorized as Stage 0/I systems by Heiderman & Evans \((2015)\), based on the criteria of HCO\(^+\) \((3-2)\) integrated line intensity of \(\geq 0.68 \text{ K km s}^{-1}\), while J182841.87 has been categorized as a Stage II system. We note that the Heiderman & Evans \((2015)\) criteria is solely based on a HCO\(^+\) detection within 14\("\) of the Spitzer c2d source position. As can be seen in Figure 6, J182902.12 show a broad HCO\(^+\) profile centered at the cloud velocity. The line model fit shows a typical Gaussian, which implies that there is no infalling envelope in this case. We have shown through SED modeling that J182902.12 is a Class Flat system and may possess a tenuous envelope. Using

**Figure 8.** SEDs with model fits for the misidentified sources. The top panel shows the two possible fits for J182841.87. Red, green, and blue lines indicate the individual contributions from the envelope, disk, and stellar components, respectively.
the “Stage” classification scheme also places J182902.12 in the Stage I-T category. Therefore, the Stage 0/I classification based on the HCO$^+$ detection is inconsistent with other schemes. For the case of J182855.78, results from line modeling indicate that the observed HCO$^+$ asymmetry can be explained by an infalling envelope model (Figure 6). The SED modeling results and the Stage classification also indicate that this is a genuine case of a Class 0/I or Stage 0/I system. For the case of J182956.67, the HCO$^+$ profile shows a hint of a blueshifted asymmetry and self-absorption at higher velocities (Figure 6). The results from line modeling for J182956.67 suggest that there may be blue-dominated infall asymmetry due to a weak envelope component with velocities between 0.5 and 1 km s$^{-1}$ and a mass infall rate of $<10^{-7}$ $M_\odot$ yr$^{-1}$. However, both the SED model fit and the Stage classification indicate that this is a Class II edge-on disk system or a Stage II source, rather than embedded Stage 0/I systems (Table 1; column 9). On the contrary, J182841.87 can be placed in the Stage I-T category and shows signatures of a tenuous envelope (Section 3.2), unlike the Stage II classification based on HCO$^+$ nondetection. The envelope mass in this system is too low to produce any detectable emission in a high-density tracer such as the HCO$^+$ (3-2) line. The submillimeter emission observed toward this source likely arises from a tenuous dusty envelope with low HCO$^+$ abundance. Overall, there appears to be less consistency in the classification determined from the physical structure of the system and that from the molecular line detection (Table 1; cols. 8, 9).

It is important to note that there are shortcomings with these various classification schemes; none of these schemes have been developed using very low-luminosity objects. Such low-luminosity Stage 0/I VLM/BDs are expected to have smaller disks and envelopes than typical protostars (e.g., White & Hillenbrand 2004), which could result in a Stage I-T or a Stage II classification. Since these classification schemes do not extend to the substellar mass regime, there may be a potential bias in classifying bona fide Stage 0/I or Class 0/I BDs as Stage II or Stage I-T objects. A revision of the criteria for the HCO$^+$ integrated line intensity of $>0.4$ K km s$^{-1}$ is also required for objects at the VLM end. For cases such as the proto-BD candidate J182852.76+002846.8 (Section 3.5), the HCO$^+$ emission is quite weak but shows a slightly blueshifted asymmetry that suggests the presence of an infalling envelope (Figure 9). As mentioned, using these various schemes to constrain the evolutionary stage of the system, we have found four out of seven bona fide detections to be misidentified objects. Assuming that the six proto-BD candidates undetected in the submillimeter continuum (Section 3.5) and the marginal or unresolved sources (Appendix) are all correctly classified, we have a $\sim$20% fraction of misidentified objects among Class 0/I/Flat VLM/BD sources in Serpens. In comparison, Heiderman & Evans (2015) estimate a $\sim$16% fraction of the full Class 0/I/Flat sample from Evans et al. (2009) in Serpens to be misidentified, and they suggest these to be Stage II disk-dominated objects. As we have shown in Section 3.4.1, at least four of the cases classified as Stage 0/I objects in Heiderman &
Evans (2015) are more evolved cases of Stage II/I-T objects. Thus in both of these studies, the estimate on the misidentified fraction is likely to be a lower limit.

4.2. Total Mass of the System

For the bona fide detections in the SCUBA-2 maps, we have derived the total (dust+gas) mass arising from the (envelope + disk) components of the system, \( M_{850}^{d+e} \), using the 850 \( \mu \)m flux density. These masses have been derived assuming a dust temperature, \( T_{\text{dust}} \), of 10 K, a gas to dust mass ratio of 100, and a dust mass opacity coefficient at 850 \( \mu \)m of 0.0175 cm\(^2\) g\(^{-1}\) (column 6, Table 1 in Ossenkopf & Henning 1994, corresponding to agglomerated dust grains with thin ice mantles at densities \( \sim 10^6 \) cm\(^{-3}\)). Table 4 lists the \( M_{850}^{d+e} \) estimates, along with the (envelope+disk) mass (dust+gas) obtained from the best model fit to the observed SED, \( M_{\text{dust}}^{\text{SED}} \).

For all cases except J182855.78, \( M_{850}^{d+e} \) is much higher than \( M_{\text{dust}}^{\text{SED}} \). This suggests that either \( T_{\text{dust}} \) is higher than the assumed value of 10 K or there may be more flux in the 850 \( \mu \)m beam width than arising from the compact source itself. We have estimated the value for \( T_{\text{dust}} \) that provides the best match between \( M_{850}^{d+e} \) and \( M_{\text{dust}}^{\text{SED}} \) (Table 4; column 4). A higher \( T_{\text{dust}} \) value of \( \sim 20-40 \) K provides a better match between the two estimates. Considering that we have used the masking technique that should minimize the effects of beam dilution (Section 2.2), the system temperature can be expected to be warmer than 10 K. The high \( T_{\text{dust}} \) values are consistent with the average dust temperatures of \( \sim 30-40 \) K found from radiative transfer models of low-mass Class I protostars (e.g., Jørgensen et al. 2011; Harsono et al. 2015). Also listed in Table 4 (columns 5, 6) are the envelope or disk size obtained from the SED fit (\( S_{\text{mod}} \)) and the spatial extent of the source as measured in the 850 \( \mu \)m image (\( S_{\text{obs}} \)). As can be seen, \( S_{\text{obs}} \) is \( \sim 2-10 \) times larger than the size \( S_{\text{mod}} \) from the SED fit. This can be expected given the large beam size in single-dish observations, and \( S_{\text{obs}} \) is likely an upper limit on the actual size of the system; interferometry can provide a more robust measurement on the source size. The larger than expected \( S_{\text{obs}} \) suggests that \( M_{850}^{d+e} \) could be underestimated and should be considered as an upper limit to the total (cold) mass in the system.

The derived values of \( M_{\text{dust}}^{\text{SED}} \) represent the available mass reservoir in the disk and envelope in the form of gas and dust. Low values of \( M_{\text{dust}}^{\text{SED}} \) as for J182902.12 and the last four objects in Table 4, imply that the mass of the central objects in these systems will not increase appreciably in the subsequent evolution. This conjecture is reinforced by the fact that they are classified as Flat objects, meaning that most of \( M_{\text{dust}}^{\text{SED}} \) is contained in the disk rather than in the envelope.

We now proceed with setting constraints on the mass of the central object using the measured bolometric luminosities and numerical simulations of stellar evolution described in Baraffe et al. (2012) and Vorobyov et al. (2016). These authors calculate stellar properties starting from a protostellar seed of 1.0 \( M_{\text{Jup}} \) and using the realistic mass-accretion rates derived from numerical hydrodynamics simulations of disk evolution. While Baraffe et al. use precalculated mass-accretion rates, Vorobyov et al. employ a fully self-consistent coupled evolution of the central object with disk hydrodynamic models. In both cases, the Lyon stellar evolution code is used to calculate the stellar properties (Chabrier & Baraffe 2000). We refer the reader to these works for a detailed description. Here, we simply provide the results of the numerical simulations by Vorobyov et al. (2016).

Figure 10 presents the \( L_{\text{int}}-M_{\text{obj}} \) diagram for 31 models of accreting proto-BDs and protostars, where \( L_{\text{int}} \) is the sum of accretion and photospheric luminosities of the central object (star or BD), and \( M_{\text{obj}} \) is the current mass of the central object. The data correspond to the Class I phase of stellar evolution. The horizontal dashed lines indicate the minimum and maximum bolometric luminosities of 0.05 \( L_{\odot} \) and 0.3 \( L_{\odot} \) in

| Object Name (SSTc2d +) | \( M_{\text{dust}}^{\text{SED}} \) (M\(_{\text{Jup}}\)) | \( M_{850}^{d+e} \) (M\(_{\text{Jup}}\)) | \( T_{\text{dust}} \) (K) | \( S_{\text{mod}} \) (") | \( S_{\text{obs}} \) (") |
|------------------------|--------------------------|--------------------------|----------------|----------------|----------------|
| J182902.12+003120.7    | 10 ± 3                   | 31 ± 4                   | 15             | 6 ± 0.4        | 16 ± 4         |
| J182855.78+002944.8     | 223 ± 10                 | 178 ± 24                 | 9              | 7 ± 0.4        | 16 ± 4         |
| J182949.57+011706.0     | 50 ± 13                  | 390 ± 40                 | 35             | 9 ± 2          | 8 ± 4          |
| J182841.87+003212.3     | 19 ± 6                   | 93 ± 10                  | 40             | 6 ± 1.5        | 16 ± 4         |
| J182956.67+011239.2     | 12 ± 4                   | 97 ± 9                   | 40             | 0.3 ± 0.08     | 8 ± 4          |
| J182902.84+003009.6     | 11 ± 3                   | 50 ± 5                   | 25             | 0.05 ± 0.01   | 12 ± 4         |
| J182955.69+011431.6     | 18 ± 5                   | 174 ± 19                 | 40             | 0.4 ± 0.08     | 16 ± 4         |
our sample (see Table 1). We note that the bolometric luminosity should be considered as an upper limit on the true internal luminosity, \( L_{\text{int}} \), of the system. The \( L_{\text{int}} \) is the luminosity of the central protostar that does not include any contribution from the external heating of the circumstellar envelope by the interstellar radiation field. As a rough estimate, we have used the linear least-squares relation between the observed flux at 70 \( \mu \)m and \( L_{\text{int}} \), obtained by Dunham et al. (2008) for a set of more than 80 embedded low-luminosity protostars. The \( L_{\text{int}} \) values are estimated to be \( \sim 70\%-80\% \) of \( L_{\text{bol}} \). Thus while \( L_{\text{bol}} \) may include a nonzero contribution from external heating sources, a major fraction of it comes from the central source.

As Figure 10 indicates, the lowest luminosities of \( L_{\text{bol}} = 0.05 L_\odot \) (J182902.12; \( L_{\text{int}} \sim 0.03 L_\odot \)) and \( L_{\text{bol}} = 0.06 L_\odot \) (J182956.67; \( L_{\text{int}} \sim 0.04 L_\odot \)) can only be attained by central objects with \( M_\text{obs} \) in the substellar mass regime. Taking into account that \( M^\text{SED}_{\text{int}} \) for these two objects is also low (\( \sim 10\%-12\% M_\text{up} \), Table 4), they will most likely form BDs. For the following three objects, J182841.87 (\( L_{\text{int}} \sim 0.13 L_\odot \)), J182902.84 (\( L_{\text{int}} \sim 0.09 L_\odot \)), and J182955.69 (\( L_{\text{int}} \sim 0.15 L_\odot \)), the mass of the central object is in the range of \( \sim 0.04-0.09 M_\odot \) (Figure 10). Considering their low gas+dust mass reservoirs, with \( M^\text{SED}_{\text{gas+dust}} \sim 0.01-0.02 M_\odot \), they may eventually end up in either the substellar or the VLM (\( \sim 0.1-0.3 M_\odot \)) categories. On the other hand, the remaining two objects, J182855.78 (\( L_{\text{int}} \sim 0.1 L_\odot \)) and J182949.57 (\( L_{\text{int}} \sim 0.14 L_\odot \)), with \( M^\text{SED}_{\text{int}} \sim 0.05-0.2 M_\odot \), can form VLM stars in the long run, even though the current masses of their central objects may still be in the substellar regime. For all of these cases, some of the envelope or disk material may be further photoevaporated or ejected due to jet or outflow activity.

5. SUMMARY

We have conducted a multiwavelength study to identify and characterize proto-BD candidates in the Serpens Main and Serp/G3–G6 clusters. Our study has revealed four good candidates, as well as a \( \sim 20\% \) fraction of VLM/BD sources misidentified as embedded YSOs. The lowest-luminosity source detected has an observed \( L_{\text{bol}} \sim 0.05 L_\odot \). We have conducted radiative transfer modeling of the observed SEDs and the HCO\(^+\) (3-2) line profile and have considered different classification schemes to understand the evolutionary stage of the system. For two candidates, there appears to be only a tenuous dusty envelope, consistent with being Stage I-T/Class Flat sources. Two other sources show more massive envelopes, indicative of being Stage 0/I systems. We have set constraints on the mass of the central object using the measured bolometric luminosities and numerical simulations of stellar evolution. Considering the available gas+dust mass reservoir and the current mass of the central source, three of these candidates are likely to evolve into BDs.

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APPENDIX

A.1. Marginal Detections

There are three sources that are detected at a marginal (\( \sim 2\sigma \)) level in the SCUBA-2 850 \( \mu \)m band but are undetected at 450 \( \mu \)m as well as in the PACS bands. These sources are listed in Table 5 with the SCUBA-2 photometry, which should be considered as upper limits. None of these objects have prior detections reported at 1.1 mm or in the MIPS 70 \( \mu \)m band. There are no HCO\(^+\) (3-2) observations available for these sources. For all of these cases, the best fits from modeling the SEDs including the PACS and 450 \( \mu \)m upper limits indicate the presence of an edge-on disk of a VLM of \( \sim 1-4 M_\text{up} \). The spectral slope of \( -0.29 \) for J183018.17+011416.9 is very close to the Class Flat/II boundary, so this is likely a disk-only source. It has been classified as M2.9 \( \pm 0.6 \) YSO (Gorlova et al. 2010).

Among the Class I sources, SSTc2d J183002.08+011359.0 lies in a nebulus region, and J182959.38+011041.1 lies in the filament seen in the southwest subcluster. It is likely that the excess emission seen in the mid-infrared for these sources is due to the surrounding cold material rather than the source itself, resulting in the steep rise in the SED observed between 3.6 and 4.5 \( \mu \)m and a high value for the spectral slope. SSTc2d J182959.38+011041.1 also has a very low observed \( L_{\text{bol}} \) of \( 0.008 L_\odot \), while the lowest-luminosity source detected in our SCUBA-2 observations has \( L_{\text{bol}} \sim 0.05 L_\odot \). In contrast, the marginal detection of the relatively more luminous Class I source SSTc2d J183002.08+011359.0 is perplexing. The near-infrared spectra for this source are dominated by excess continuum emission (Figure 2) and are similar to the spectra observed in extinguished VLM protostars (e.g., Connelley &
Greene 2010; Gorlova et al. 2010). Deeper scans can help further characterize these sources.

A.2. Unresolved or Confused Sources

We have found four sources that are unresolved from a nearby bright YSO in the SCUBA-2 maps. The lowest-luminosity case among these is SSTc2d J182909.05+003128.0 ($L_{\text{bol}} = 0.0078 \, L_{\odot}$), which lies at a ~2" separation from the protostar SSTc2d 182909.07+003132.4 ($L_{\text{bol}} = 1.1 \, L_{\odot}$). The two objects are also unresolved in the PACS bands. While there are separate HCO$^+$ (3-2) pointings for the two sources, the observed profiles are very similar. The HCO$^+$ profiles for both sources show weak emission centered at the systemic velocity. Considering the large beam size for these molecular line observations, the two pointings must be tracing the composite source. For another very low-luminosity case SSTc2d J182952.21+011559.1 ($L_{\text{bol}} = 0.24 \, L_{\odot}$), there is no object seen at the source location in the SCUBA-2 450 $\mu$m map. This object is at a ~13" separation from the bright protostar SSTc2d J182952.08+011547.8 and is within the beam size of ~14′/5 in the 850 $\mu$m band. This is a case of both a nondetection and an unresolved object due to the poor sensitivity and angular resolution in the SCUBA-2 maps.

Among the more luminous ($L_{\text{bol}} \sim 0.1 \, L_{\odot}$) Class Flat sources, SSTc2d J182957.66+011304.6 is in close proximity (~8") to the prestellar core JCMTSF J182956.6+011309, while the faint object SSTc2d J182949.69+011456.8 is missed in the bright halo of the prestellar core JCMTSF J182949.8+011515. The HCO$^+$ line for both sources shows strong emission with a blue-dominated axisymmetric profile, indicative of an infalling envelope. The strong emission, however, is caused by the massive sources within the 25" beam size of the HCO$^+$ observations. Due to the HCO$^+$ detection, all of these faint unresolved or confused sources have been categorized as Stage 0/1 embedded objects (Heiderman & Evans 2015), which is incorrect. The classification for these objects is questionable as the sharp rise in the SED could be contaminated by the nearby protostar.

An additional source, SSTc2d J183000.30+010944.7 ($\alpha_{\text{IR}} = -0.12, \, L_{\text{bol}} = 0.09 \, L_{\odot}$), is undetected in our continuum observations. This object is located in the filament at the very south of the southwest subcluster. There is no point source seen in the SCUBA-2 and PACS maps at the object location except confusion noise, which may have caused the Class I SED shape and the high $L_{\text{bol}}$ for this object. The actual source might be fainter than the observed luminosity. The HCO$^+$ line shows a symmetric profile centered at the systemic velocity. The $\alpha_{\text{IR}}$ for this source is closer to the Class Flat/Class II threshold, and this might be a Class II object instead of the Stage 0/1 classification determined from HCO$^+$ detection. Higher-resolution observations are required to resolve these sources and confirm their evolutionary stage.

### Table 5

| Object Name (SSTc2d | R.A. [J2000] | Decl. [J2000] | $\alpha_{\text{IR}}$ | $L_{\text{bol}}$ ($L_{\odot}$) | $T_{\text{bol}}$ (K) | Class* | SCUBA-2* |
|--------------------|-------------|---------------|-----------------|----------------|-------------|--------|---------|
| J183002.08+011359.0 | 18:30:02.08 | +01:13:59.0 | 0.36 | 0.09 | 510 | I, I | <54 | <96 |
| J182959.38+010041.1 | 18:29:59.38 | +01:10:41.1 | 0.51 | 0.008 | 440 | I, I | <17 | <57 |
| J183018.17+011416.9 | 18:30:18.17 | +01:14:16.9 | −0.29 | 0.13 | 1200 | Flat, II | <15 | <41 |
| J182952.21+011559.1 | 18:29:52.21 | +01:15:59.1 | 1.87 | 0.024 | 230 | I, I | <273 | <137 |

**Notes.**

* The estimates for $\alpha_{\text{IR}}$, Class, $L_{\text{bol}}$, and $T_{\text{bol}}$ are from Evans et al. (2009).

* First value is based on $\alpha_{\text{IR}}$ and second value is based on the $T_{\text{bol}}$ classification ranges, from Evans et al. (2009).

* Upper limits.

### A.3. Consequences of Uncertainties

For three sources in our study, J182902.12, J182855.78, and J182902.84, there is 1.1 mm photometry available from CSO Bolocam observations presented in Enoch et al. (2009). For all three cases, the 1.1 mm flux density is higher than the 850 $\mu$m photometry and is notably offset from the best model fit to the SED (Figures 5, 8). The beam size for the CSO Bolocam is almost twice that of the SCUBA-2 850 $\mu$m band (~14′/5). As can be seen in the 850 $\mu$m images, there is a bright source about 18″ away from the target J182855.78, while J182902.84 and J182902.12 lie in regions where diffuse nebulosity is seen about 20″–25″ from the target position (Figures 4, 7). It appears that, for all three cases, the flux within the Bolocam 30″ beam is likely enhanced due to the contribution from another source within the beam or a higher contribution from the surrounding material, whereas there is less contribution from these contaminants in the 14″ SCUBA-2 beam size.

As discussed in Section 2.2, we have applied the masking technique to the SCUBA-2 maps, which can provide the optimum source contribution, particularly for such faint objects. The 850 $\mu$m photometry is thus more reliable, since a similar masking technique has not been applied to the published 1.1 mm photometry. We have opted not to use the 1.1 mm point in fitting the SEDs; any model that provides a good fit to this point misses the submillimeter and far-infrared points and is not a good fit to the overall SED. Nevertheless, we have attempted to fit the 1.1 mm point for the case of J182902.12, which shows the largest offset in the 1.1 mm photometry (Figure 5). The model that provides a good fit to the 1.1 mm point shows more than half an order of magnitude higher fluxes at far-infrared and submillimeter wavelengths than the observed photometry and is for an envelope mass of ~14 $M_{\text{Jup}}$ and a disk mass of ~3 $M_{\text{Jup}}$, higher than the mass estimates obtained from the best fit (Table 3). Likewise, the envelope+disk mass for the J182855.78 system from a model that fits the 1.1 mm point is ~260 $M_{\text{Jup}}$ while a similar fit to J182902.84 is for a disk mass of ~24 $M_{\text{Jup}}$. As discussed in Section 4.2, $M_{\text{obj}}$ for J182902.12 is in the substellar regime, and with an addition of envelope+disk mass of ~18 $M_{\text{Jup}}$, this system is still likely to reach a final mass below the substellar limit, while J182855.78 and J182902.84 will likely evolve into VLM stars. However, the reduced-$\chi^2$ value for such a fit is >5, much poorer than the fit obtained excluding the 1.1 mm point.

### Table 5

**Marginal Detections in the SCUBA-2 Maps**

| Object Name | R.A. [J2000] | Decl. [J2000] | $\alpha_{\text{IR}}$ | $L_{\text{bol}}$ ($L_{\odot}$) | $T_{\text{bol}}$ (K) | Class* | SCUBA-2* |
|-------------|-------------|---------------|-----------------|----------------|-------------|--------|---------|
| J183002.08+011359.0 | 18:30:02.08 | +01:13:59.0 | 0.36 | 0.09 | 510 | I, I | <54 | <96 |
| J182959.38+010041.1 | 18:29:59.38 | +01:10:41.1 | 0.51 | 0.008 | 440 | I, I | <17 | <57 |
| J183018.17+011416.9 | 18:30:18.17 | +01:14:16.9 | −0.29 | 0.13 | 1200 | Flat, II | <15 | <41 |
| J182952.21+011559.1 | 18:29:52.21 | +01:15:59.1 | 1.87 | 0.024 | 230 | I, I | <273 | <137 |
The alternative total mass obtained by fitting the 1.1 mm photometry is thus not reliable. Furthermore, in order to compare the total masses and physical properties for all sources in this study (Section 4), it is important to be consistent and use the measurements obtained by applying the same methodologies on the same set of data, rather than using the masses obtained from the 1.1 mm photometry for just three sources.

Finally, the distance to Serpens is still uncertain, as discussed in Section 1. While nearly all distance estimates are within ~260–310 pc, the Dzib et al. (2010) measurement of 415 pc would imply that the $L_{bol}$ estimates for our targets would be higher by a factor of ~2.5. Based on the arguments provided in Section 4.2, this would place the lowest-luminosity sources J182902.12 and J182956.67 in the VLM regime, while the rest will likely evolve into low-mass stars.

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