STRONG GRAVITATIONAL LENSING BY THE SUPER-MASSIVE cD GALAXY IN ABEll 3827
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Abstract

We have discovered strong gravitational lensing features in the core of the nearby cluster Abell 3827 by analyzing Gemini South GMOS images. The most prominent strong lensing feature is a highly magnified, ring-shaped configuration of four images around the central cD galaxy. GMOS spectroscopic analysis puts this source at $z \sim 0.2$. Located $\sim 20''$ away from the central galaxy is a secondary tangential arc feature which has been identified as a background galaxy with $z \sim 0.4$. We have modeled the gravitational potential of the cluster core, taking into account the mass from the cluster, the brightest cluster galaxy (BCG), and other galaxies. We derive a total mass of $(2.7 \pm 0.4) \times 10^{15} M_\odot$ within $37 h^{-1}$ kpc. This mass is an order of magnitude larger than that derived from X-ray observations. The total mass derived from lensing data suggests that the BCG in this cluster is perhaps the most massive galaxy in the nearby universe.

Key words: galaxies: clusters: individual (Abell 3827) – galaxies: elliptical and lenticular, cD – galaxies: evolution – galaxies: formation – gravitational lensing: strong

1. INTRODUCTION

As the densest galaxy environments known, the cores of massive clusters are expected to host the strongest dynamical evolution. Such cluster cores are found to be dominated by early-type D or cD galaxies that are also the brightest cluster galaxies (BCGs). These have extended luminous haloes (Schombert 1987), which are not necessarily smooth (Johnstone et al. 1991) and contain multiple or complex nuclei (Rood & Leir 1979). They are located close to the peak of the X-ray emission (Jones & Forman 1984) and near the kinematical centers of their clusters (Quinta & Lawrie 1982; Quintana et al. 2000).

There has been a long-running debate over the extent to which BCGs can be assembled through continuous merging of galaxies in the cluster potential (“galactic cannibalism”) versus early merging during cluster collapse (e.g., West 1994; Dubinski 1998). However, recent work has shown that BCGs are likely to form via “dry” or dissipationless major mergers of smaller early-type galaxies (e.g., van Dokkum 2005; Bell et al. 2006; de Lucia & Blaizot 2007; Whitaker & van Dokkum 2008). These dry mergers are thought to be the primary mechanism through which massive galaxies grow from $z \sim 1$ to the present day and continue populating the upper end of the mass function without changing the overall mass density of elliptical systems. It has been pointed out that the merger rate increases both with the stellar mass of galaxies and with age (Liu et al. 2009). Therefore, there is a large fraction of major dry mergers in the nearby universe. Indeed, as many as 3.5% of BCGs show ongoing evidence of mergers (Liu et al. 2009), and about 10%–20% of massive galaxies have undergone a dry merger in the last gigayear (Khochfar & Silk 2009). Such cluster cores are found to be dominated by early-type D or cD galaxies that are also the brightest cluster galaxies (BCGs). These have extended luminous haloes (Schombert 1987), which are not necessarily smooth (Johnstone et al. 1991) and contain multiple or complex nuclei (Rood & Leir 1979). They are located close to the peak of the X-ray emission (Jones & Forman 1984) and near the kinematical centers of their clusters (Quinta & Lawrie 1982; Quintana et al. 2000).

2. OBSERVATIONS

All imaging and spectroscopic data were collected with the Gemini Multi-Object Spectrograph (GMOS; Hook et al. 2004) at the Gemini South telescope, in queue mode.

The cluster was imaged with the $g'$, $r'$, and $i'$ filters on 2007 November 7 (UT), in photometric conditions, with good seeing (0.4–0.6'). Observations were processed with the Gemini IRAF package (version 1.8) in the standard manner. Calibration on the standard magnitude system was achieved using observations of stars from Landolt (1992). Object detection and photometry were performed on the $r'$-band image (2007 November 7 UT) with the program SExtractor (Bertin & Arnouts 1996). The MAG_AUTO parameter was adopted as the total magnitude for the detected objects. Colors were derived by measuring fluxes inside a fixed circular aperture of 1.5'' (~10 pixels) in all filters, corresponding to a physical aperture of 2.75 $h^{-1}$ kpc at the core of Abell 3827.
Figure 1. Color-composite image of Abell 3827. The field of view is 5′ × 5′ (≈0.55 × 0.55 h⁻² Mpc² at the distance of the cluster). The cD galaxy and its asymmetric halo are clearly seen at the center of the image. Top left panel: the field of view is approximately 40′ × 40′. The five nuclei contained in the center are labeled N.1 to N.5 and two stars are labeled S. Bottom right panel: the field of view and orientation are similar to that of the top left panel. The “central” arcs are labeled A.1 to A.4 and the tangential arc is labeled B.1. With the central nuclei removed and the image display scale chosen to maximize contrast, all of the gravitational lensing features are more distinct.

distance of the cluster. All objects with SExtractor stellarity flag \( \leq 0.8 \) were selected as galaxies. We estimate that the catalog is complete down to \( r' = 23 \) mag, since the number counts start to turn over at this value. The final catalog contains the total magnitudes, colors, and structural parameters for 747 galaxies brighter than \( r' = 23 \) mag. Figure 1 shows the color-composite \( g', r', i' \) image of the central region of Abell 3827 (0.55 × 0.55 h⁻² Mpc²). Five central elliptical galaxies (N.1 to N.5) are embedded in a common asymmetric halo (top left panel). The asymmetric halo has an ellipsoidal form, with \( a \sim 1'.3 \) and \( b \sim 1' \) for the semimajor and semiminor axes, respectively. Strong lensing features detected around the central cD galaxy (at radii \( \sim 8'' \) and \( \sim 20'' \)) are visible in the images located in the top left and bottom right panels of Figure 1.

Long-slit spectra of the central five elliptical galaxies were obtained on 2007 Jun 3 (UT) during bright time, through thin-cloud cover and \( \sim 1'' \) seeing. The observations were performed using the nod-and-shuffle mode, with the R400 grating centered at 8100 Å; this setup was chosen to minimize the effect of moon illumination and to obtain good sky subtraction. To avoid second-order contamination from blue light in the red part of these spectra, the OG515 blocking filter was used. In addition, multi-object spectroscopic (MOS) observations of galaxies in the Abell 3827 field were obtained on 2007 September 7 (UT), during dark time, in photometric conditions, and with seeing between 0.7' and 0.9'. Two masks were observed with 0.75 slit widths, using the B600 grating centered at 5220 Å. During the observations, all spectra were dithered to account for the gaps in the GMOS CCDs. Spectra were reduced and calibrated in the standard manner for each mode.

3. RESULTS

We have measured radial velocities for 67 galaxies using cross-correlation techniques or emission-line fitting. Of these 67 galaxies, 56 are within \( \pm 2500 \) km s⁻¹ of the cluster redshift. Using the bi-weight estimator for location and scale (Beers et al. 1990), we calculate an average velocity for the cluster of \( 29663 \pm 155 \) km s⁻¹ and a line-of-sight velocity dispersion of \( 1142 \pm 125 \) km s⁻¹, with 55 member galaxies. The virial mass inside a radius of 0.3 h⁻¹ Mpc is \( (3.86^{+0.46}_{-0.35}) \times 10^{14} \) h⁻¹ M☉, with uncertainties at the 68% confidence intervals.

In the cluster core, we confirm membership for four of the central elliptical galaxies. There are also two stars in the foreground (see Figure 1). Three of the core galaxies have similar colors (see Table 1) and their velocities are consistent with the average cluster velocity (within \( \sim 300 \) km s⁻¹). Galaxy N.4 is slightly bluer than N.1, N.2, and N.3 and has a measured
radial velocity $\sim 1000$ km s$^{-1}$ higher than the systemic velocity. Galaxy N.5 may not be in the cluster, but the background.

### 3.1. Strongly Lensed Features in the Cluster Core

Thanks to the superb image quality delivered by GMOS and the Gemini South telescope, we have been able to detect several strongly lensed features around the central cD galaxy in the core of the cluster. A thin, extended tangential arc (B.1 in Figure 1) is seen at $\sim 20''$ southeast from the central galaxy, with apparent length $\sim 11''$ and thickness $\sim 0.7''$. Surrounding the central core galaxies is a second and more prominent arc feature—a highly magnified, ring-shaped configuration of four images around the central cD galaxy, with a radius of $\sim 8''$. Only three of the arc components are visible in the inset color image in Figure 1 (A.1, A.2, and A.3). The fourth image (A.4 in Figure 1) is between galaxies N.1 and N.2 and is obscured by the halo of the cD galaxy. All four features can be seen in detail in the bottom right panel of Figure 1 (see the caption for details). These features are very reminiscent of the lensed images seen in Cl0024+1654 (Colley et al. 1996), although the source galaxy in the present case is closer than $z \sim 1$ (see below).

The redshifts of the lensed sources were determined from long-slit observations taken on 2007 November 15 and 16, using the R150 grating centered at 7150 Å. For the feature with the ring-shaped configuration (system A), we derived a redshift for one of the images located northwest of the cluster core (A.2). Two emission lines (H$\alpha$ and [N$\text{II}$]$\lambda6583$ Å) are seen in the spectrum for A.2; unfortunately no other lines were detected, due to light contamination from the cD halo. From the two emission lines, the estimated redshift for A.2 is $z = 0.20443 \pm 0.00073$.

The symmetric distribution of images in system A and the ring-shaped structure connecting arcs A.1, A.2, and A.3 (and possibly A.4) suggest that the images are of the same source galaxy. Therefore, we assume that all four images come from the same source and have the same redshift. This assumption should be taken with caution, without spectroscopy of the remaining system A images for confirmation, but is supported by our modeling, discussed in the next section.

The subtraction of light from the halo of the cD galaxy, as well as subtraction of all galaxies and the stars from the cluster core, highlights the ring-shaped structure connecting arcs A.1, A.2, A.3, and possibly A.4 (see Figure 1). The existence of this structure only $\sim 15h^{-1}$ kpc from the center of the cD galaxy gives us a unique opportunity to study the dynamics and mass distribution of a large BCG at unprecedentedly small scales.

Using several emission lines ([O$\text{II}$] $\lambda3727$ Å, H$\beta$, [O$\text{II}$] $\lambda\lambda4959, 5007$ Å) seen in the spectrum of the southeast tangential arc, B.1, we derived a redshift of $z = 0.40825 \pm 0.00072$. The relatively long, thin appearance of this arc suggests that the source galaxy lies very close to the inner fold caustic, assuming the lens cluster has an elliptical spatial distribution of mass (e.g., see Fort & Mellier 1994; Narayan & Bartelmann 1999). If true, there must be three fainter counterimages: one to the northeast, one to the southwest, and one at the center of the cluster core.

### 3.2. Strong Lensing Modeling and Mass Determination

To produce the features described in the previous section, the cluster core must be very massive. In order to reconstruct the mass distribution in the core (the region enclosed by system A), we used the parametric model implemented in the publicly available LENSTOOL$^4$ ray-tracing code (Jullo et al. 2007). To model the lens, we used two clumps of different scales—a large-scale halo, representing both the matter inside the giant cD galaxy and the dark matter halo of the cluster, and four smaller-scale clumps, modeling the small-scale perturbations associated with the galaxies embedded in the cD galaxy. For all components we adopt a dual Pseudo-Isothermal Elliptical Mass Distribution (dPIEMD; see Elíasdóttir et al. 2007). The dPIEMD can be characterized by seven parameters: the center position $(X, Y)$, the ellipticity $\epsilon$, the position angle $\theta$, and the parameters of the density profile: the velocity dispersion $\sigma_0$ and two characteristic radii $r_{\text{core}}$ and $r_{\text{cut}}$. For the large halo, we left most parameters free, with broad uniform priors. Only $r_{\text{cut}}$ remained fixed at a value of $1500h^{-1}$ kpc. In the small-scale clumps, the parameters of the density profiles were scaled as a function of their galaxy luminosities (see Jullo et al. 2007, and references therein). Using as a scaling factor the luminosity $L^*$, associated with the $g$-magnitude of central galaxy N.3 (see Figure 1 and Table 1), we searched for the values of $\sigma_0^*$ and $r_{\text{cut}}^*$ that yield the best fit, fixing $r_{\text{core}}^*$ at $0.15h^{-1}$ kpc. The best-fitting parameters of the model are listed in Table 2.

Because the ring structure shows multiple subcomponents, we selected the most reliable set of multiple images and associated them with three different background sources in order to make our fit. Considering an uncertainty in the position of any image equal to $0''.2$, we found after the optimization in the image plane $\chi^2_{\text{DOF}} = \chi^2/\text{DOF} = 69/4 \sim 17$. Our model reproduces well the positions of the observed subcomponents, with a mean scatter less than $0''.8$. To highlight this, we show in Figure 2, the

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**Table 1**

Main Data of the Five Central Ellipticals Galaxies Located in the Core of Abell 3827

| Galaxy | ID | R.A.(2000) (h m s) | Decl.(2000) (° ′ ″) | $V_{\text{hel}}$ (km s$^{-1}$) | $\delta V_{\text{hel}}$ (km s$^{-1}$) | $r'$ (arcsec) | $i'$ (arcsec) | $r'$ (mag) | $i'$ (mag) | $(r'-i')_{\text{N3}}$ (mag) |
|--------|----|-------------------|--------------------|-----------------|----------------------|-------------|-------------|-------------|-------------|------------------|
| N.1    | 1232 | 22:01:53.99       | $-59:56:45.14$     | 29548           | $-115$               | 13.88       | $-1.92$     | 17.03       | 16.51       | 0.51             |
| N.2    | 1234 | 22:01:53.31       | $-59:56:43.14$     | 29763           | 100                   | 3.75        | 0.08       | 16.19       | 15.67       | 0.52             |
| N.3    | 1235 | 22:01:52.69       | $-59:56:42.11$     | 29951           | 288                   | $-5.75$     | 2.00       | 16.18       | 15.56       | 0.55             |
| N.4    | 1236 | 22:01:51.76       | $-59:56:46.12$     | 28590           | $-1073$               | $-4.63$     | $-2.91$    | 16.23       | 15.81       | 0.43             |
| N.5    | 1237 | 22:01:52.79       | $-59:56:37.73$     | 33817           | 4154                  | $-4.13$     | 5.48       | 18.08       | 17.70       | 0.35             |

**Notes.** (1) Galaxy number in the Source Extractor catalog, (2) galaxy ID in Figure 1, (3) and (4) R.A. and decl. (J2000.0), (5) heliocentric radial velocity, (6) difference between the heliocentric radial velocities of each component and the systemic velocity, (7) and (8) coordinates relative to the center of mass of the system located at $\alpha_{2000} = 22^h01^m53^s06$, $\delta_{2000} = -59^\circ43'43''/22$, (9) and (10) total magnitudes, and (11) galaxy colors measured inside a fixed aperture of diameter $1''$.5.
produced by strong galaxy–galaxy lensing events. With these three points. For clarity, we do not depict the demagnified images show with small asterisks (cyan) the positions in the source plane associated the measured image positions. Green crosses represent the predicted image caustic line (gray) for a source at North is up and east is left. The external critical line (red) and the associated 

It is important to note that our model is oversimplified, since we are assuming spherical halos for the galaxies. A detailed model of Abell 3827 is beyond the scope of the present work. Given our best model (see Table 2), we calculated the total mass inside a radius of \(20^\prime\) (location of the tangential arc B.1 at \(\sim 37 \ h^{-1} \) kpc) and found \(M = (2.7 \pm 0.4) \times 10^{13} \ M_\odot\). This mass is slightly greater than that enclosed in Abell 1689 within a similar radius (see Figure 6 in Limousin et al. 2007). We estimated the mass of the cD galaxy as being \(\sim 30\%–50\%\) of the total mass of the cluster within \(20^\prime\) (Limousin et al. 2007). This implies a cD galaxy mass between \(8 \times 10^{10} \ M_\odot\) and \(1.3 \times 10^{13} \ M_\odot\). Even assuming a conservative value of \(30\%\), this would mean that the central cD galaxy in Abell 3827 is perhaps the most massive galaxy observed in the local universe.

Table 2

| Comp.                  | X (\(^\circ\)) | Y (\(^\circ\)) | \(\epsilon\) | \(\theta\) (kpc) | \(r_\text{core}\) (kpc) | \(r_\text{cut}\) (kpc) | \(\sigma_0\) (km s\(^{-1}\)) |
|------------------------|---------------|---------------|--------------|-----------------|-----------------|-----------------|----------------|
| Cluster+cD             | -0.9          | -2.5          | 0.1          | 56              | [1500]          | [0]              | 1174           |
| N.3(L\(^+\))           | [2.77]        | [2.82]        | 0.0          | [0]             | [0.15]          | 8               | 256            |

Notes. Coordinates are relative to the center of image 2 (\(\alpha_\text{2000} = 22^h01^m53.15^s\), \(\delta_\text{2000} = -59^\circ56'43.31''\)). Values in square brackets are not optimized.

Using the strong-lensing model presented above, we have analyzed the mass distribution to unprecedented spatial resolution. The cluster core is spatially very concentrated. Even if we assume a conservative fraction for the mass of the central cD galaxy (see above), this galaxy is very massive. The cD galaxy could be an extreme example of the effects of dry mergers on the mass of BCGs; dry mergers can produce an increase (by a factor of up to 3) in the dark-matter-to-stellar mass ratio for the most massive systems at present (Ruszkowski & Springel 2009).

We have compared the total mass derived from our strong-lensing analysis with that derived from the X-ray gas. We have calculated the total mass as a function of radius for Abell 3827, based on a recent X-ray analysis by L. Valkonen et al. (2010, private communication). We derive the total cluster mass by assuming hydrostatic equilibrium and using their best fit for the three-dimensional distribution of X-ray emitting gas. L. Valkonen et al. (2010, private communication) report \(M_{200} = 9.7 \times 10^{14} \ M_\odot\) inside a radius \(r_{200} = 2.84 \ h^{-1} \) Mpc, derived from the \(M–T\) relation by Salhén et al. (2009) for a cluster temperature of \(kT = (7.15 \pm 0.18)\) keV. Figure 3 shows the radial distribution of the total mass derived from the X-ray gas and from the strong-lensing analysis. For the X-ray gas, we plot mass profiles (1) for a typical cluster with \(r_\text{core} = 250 \ h^{-1} \) kpc and \(\beta = 0.7\) (dash-dotted line) and (2) for the derived total mass of Abell 3827 within 37 \(h^{-1}\) kpc (solid vertical line). Finally, we plot the mass profile for Abell 3827 derived from strong lensing as a dashed line. At all radii, the mass derived from strong lensing is at least a factor of 10 larger than that derived from X-ray data. Discrepancies between strong-lensing cluster masses and X-ray cluster masses for the same clusters have been reported previously in the literature (e.g., Gitti et al. 2007; Halkola et al. 2008). Moreover, significant biases of up to \(\sim 50\%\) may be introduced into strong-lensing mass estimates when models are extrapolated outside the Einstein ring (Meneghetti et al. 2010). Analysis using both strong-lensing and X-ray methods at small radii shows that differences can grow by as much as an order of magnitude (Verdugo et al. 2007).

How can we explain this discrepancy? For the lower limit on mass, within 3\(\sigma\) error of the strong-lensing fit, we obtain \(6.3 \times 10^{12} \ M_\odot\) inside the ring structure of radius \(10^\prime\) (\(\sim 19 \ h^{-1} \) kpc). Hence, the most likely uncertainties, related to extrapolation of the strong-lensing model to larger radii, are unable to
explain our factor-of-ten discrepancy. Assuming a spherical mass distribution, the mass inside system A is \((9.6 \pm 0.9) \times 10^{12} M_\odot\). The large difference between the strong-lensing mass and X-ray mass could be related to the dynamical state of the cluster, to the mass inferred from the X-ray gas, or to both effects. Total X-ray cluster masses derived from flat-core density profiles may be underestimated by at least a factor of 2 within the central \(\sim 30 h^{-1}\) kpc (Voigt & Fabian 2006). Simulations (Meneghetti et al. 2010) of strong-lensing models have shown that physical substructures along the line of sight to the cluster can overestimate the derived total mass by a factor of 2; in fact, our GMOS spectra reveal a bi-modality in the velocity distribution of cluster galaxies, suggesting that Abell 3827 is presently merging. In all, these effects can produce a factor of about 4, still insufficient to explain the large discrepancy between strong-lensing and X-ray derived masses. Moreover, if we take out the factor of 4 due to the two effects described above, we are left with a cD galaxy with a mass of \(\sim 2 \times 10^{12} M_\odot\), which appears to be one of the most massive galaxies known in the local universe. Detailed studies of the central objects at different wavelengths will reveal insights into the origins and subsequent evolution of this peculiarly massive galaxy.

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