Search for a New Heavy Gauge Boson \( W' \) with Electron+\( E_T \) Event Signature in \( pp \) collisions at \( \sqrt{s} = 1.96 \) TeV

T. Aaltonen,21 B. Álvarez González,9 S. Amerio,41 D. Amidei,32 A. Anastassov,36 A. Annovi,17 J. Antos,12 G. Apollinari,15 J.A. Appel,15 A. Apresyan,46 T. Arisawa,56 A. Artikov,13 J. Asaadi,51 W. Ashman,15 B. Auerbach,59 A. Aurisano,51 F. Azfar,40 W. Badgett,15 A. Barbaro-Galtieri,26 V.E. Barnes,46 B.A. Barnett,23 P. Barria,44 P. Bartos,12 M. Bauer,41 G. Bauer,30 F. Bedeschi,44 D. Beecher,28 S. Behari,23 G. Bellentini,44 J. Bellinger,58 D. Benjamin,14 A. Beretvas,15 A. Bhatti,30 M. Binkley,15 D. Bisello,41 I. Bizjak,29 K.R. Bland,5 B. Blumenfeld,23 A. Bocci,14 A. Bodek,37 D. Bortoletto,46 J. Boudreau,45 A. Boveia,11 B. Brau,15 L. Brigliadori,26 A. Brisida,12 C. Bromberg,33 E. Brucken,21 M. Bucciantonio,64 J. Budagov,13 H.S. Budd,47 S. Budd,22 K. Burkett,15 G. Busetto,41 P. Bussey,19 A. Buzatu,31 C. Calancha,29 S. Camarda,4 M. Campanelli,33 M. Campbell,32 F. Canelli,12 A. Canepa,43 B. Carls,22 D. Carlsmith,58 R. Carosi,44 S. Carrillo,6 S. Carron,15 B. Casal,15 M. Casarsa,15 A. Castro,6 P. Catastini,15 D. Cauz,52 V. Cavaliere,44 M. Cavalli-Sforza,4 A. Cerri,26 L. Cerrito,29 Y.C. Chen,7 M. Chertok,7 G. Chlachidze,15 F. Chlebana,15 K. Cho,25 D. Chokheli,13 J.P. Chou,29 W.H. Chung,58 Y.S. Chung,47 C.I. Ciobanu,42 M.A. Ciocci,44 A. Clark,18 G. Compostella,41 M.E. Convery,15 J. Conway,7 M. Corbo,42 M. Cordelli,17 C.A. Cox,7 D.J. Cox,7 F. Crescioli,44 C. Cueva Almenar,59 J. Cuevas,9 R. Culbertson,15 D. Dagenhart,15 N. d’Ascanzo,42 M. Datta,15 P. de Barbaro,47 S. De Cecco,49 G. De Lorenzo,4 M. Dell’Orso,44 C. Deluca,4 L. Demortier,48 J. Deng,14 M. Denino,6 F. Devoti,21 M. d’Errico,41 A. Di Canto,64 B. Di Ruzza,44 J.R. Dittmann,5 M. D’Onofrio,27 S. Donati,26 P. Dong,15 T. Dorigo,12 K. Ebina,56 A. Eppig,32 R. Erbacher,7 J. Errede,22 N. Ershaidat,9,42 R. Eusebi,51 H.C. Fang,26 S. Farrington,40 M. Feindt,24 J.P. Fernandez,29 C. Ferrazza,44 R. Field,16 G. Flanagan,46 R. Forrest,7 M.J. Frank,5 M. Franklin,20 J.C. Freeman,15 I. Furic,16 M. Gallinaro,46 J. Galyardt,10 J.E. Garcia,18 A.F. Garfinkel,46 P. Garosi,44 H. Gerberich,22 E. Gerchtein,15 S. Giaguri,49 V. Giakoumopoulos,3 P. Giannetti,44 K. Gibson,45 C.M. Ginsburg,15 N. Giokaris,3 P. Giromini,17 M. Giunta,44 G. Giurgiu,23 V. Glagolev,13 D. Gleeninski,15 M. Gold,35 D. Goldin,51 N. Goldschmidt,16 A. Golossanov,15 G. Gomez,9 G. Gomez-Ceballos,30 M. Goncharov,15 O. González,29 I. Gorelov,35 A.T. Goshaw,14 K. Goulianos,48 A. Gresele,41 S. Grinstein,4 C. Grosso-Pilcher,13 R.C. Group,55 J. Guimaraes de Costa,20 Z. Gunay-Unalan,33 C. Haber,26 S.R. Hahn,15 E. Halkiadakis,50 A. Hamaguchi,39 J.Y. Han,47 F. Happacher,17 K. Harra,53 D. Hare,50 M. Hare,54 R.F. Harr,57 K. Hatakeyama,5 C. Hayes,40 M. Heck,24 J. Heinrich,43 M. Herndon,58 S. Hewamanage,5 D. Hidas,50 A. Hocker,15 W. Hopkins,15 D. Horn,24 S. Hou,1 R.E. Hughes,37 M. Hurwitz,11 U. Husemann,59 N. Hussain,31 M. Hussein,31 J. Huston,33 G. Intorzi,44 M. Iori,49 A. Ivanov,7 E. James,15 D. Jang,10 B. Jayatilaka,4 E.J. Jeon,25 M.K. Jha,6 S. Jindariani,15 W. Johnson,46 M. Jones,46 K.K. Joo,25 S.Y. Jun,10 T.R. Junk,15 T. Kamon,51 P.E. Karchin,57 Y. Kato,39 W. Ketchum,11 J. Keung,43 V. Kholytovich,51 B. Kilminster,15 D.H. Kim,25 H.S. Kim,25 H.W. Kim,25 J.E. Kim,25 M.J. Kim,17 S.B. Kim,25 S.H. Kim,53 J.K. Kim,11 N. Kimura,56 M. Kirby,15 K. Klimenko,16 K. Kondo,56 D.J. Kong,25 J. Konigsberg,16 A.V. Kotwal,14 M. Kreps,24 J. Kroll,43 D. Krop,11 N. Krumnack,5 M. Kruse,14 V. Krutelyov,51 T. Kuhr,24 M. Kurata,53 S. Kwang,11 A.T. Laasanen,46 S. Lami,44 S. Lamml,15 M. Lancaster,28 R.L. Lander,7 K. Lannon,37 A. Lath,50 G. Latino,44 I. Lazzizzera,11 T. LeCompte,2 E. Lee,51 H.S. Lee,11 J.S. Lee,25 S.W. Lee,51 S. Lee,66 S. Leone,44 J.D. Lewis,15 C.-J. Lin,26 J. Linacre,40 M. Lindgren,15 E. Lipeles,43 A. Lister,18 D.O. Litvintsev,15 C. Liu,45 Q. Liu,46 T. Liu,15 S. Lockwitz,59 N.S. Locke,43 A. Loginov,59 D. Lucchesi,44 J. Lucek,24 P. Lujan,26 P. Lukens,15 G. Lungu,48 J. Lys,46 R. Lysak,21 R. Madrak,15 K. Maeshima,15 K. Makhoul,30 P. Maksimovic,23 S. Malik,48 G. Manca,26 A. Manousakis-Katsikakis,5 F. Margaroli,46 C. Marino,24 M. Martinez,4 R. Martinez-Ballarin,29 P. Mastrandrea,49 M. Mathis,23 M.E. Mattsson,57 P. Mazzanti,6 K.S. McFarland,47 P. McIntyre,51 R. McNulty,27 A. Mehta,27 P. Mehtala,21 A. Menzone,44 C. Mesropian,48 T. Miao,15 D. Mietlicki,32 A. Mitra,1 H. Miyake,53 S. Moed,20 N. Mogg,6 M.N. Mondragon,15 C.S. Moon,25 R. Moore,15 M.J. Morello,15 J. Morlock,24 P. Movilla Fernandez,15 A. Mukherjee,7 Th. Muller,24 P. Murat,15 M. Musiela,6 J. Nachtman,15 Y. Nagai,53 J. Nagano,56 I. Nakano,38 A. Napier,54 J. Nett,58 C. Neu,55 N.S. Neuberger,22 J. Nielsen,26 L. Nodulman,2 O. Norniella,22 E. Nurse,28 L. Oakes,40 S.H. Oh,14 Y.D. Oh,25 I. Oksuzian,39 T. Okusawa,39 R. Orava,21 L. Ortolan,4 S. Pagan Griso,40 C. Pagliarone,52 E. Palencia,9 V. Papadimitriou,15 A.A. Paramonov,2 J. Patrick,15 G. Pauletta,15 M. Paulini,10 C. Paus,30 D.E. Pellett,7 A. Penzo,52 T.J. Phillips,14 G. Piacentino,44 E. Pianori,43 J. Pilot,37 K. Pitts,22
We present a search for a new heavy charged vector boson $W'$ decaying to an electron-neutrino pair in $p\bar{p}$ collisions at a center-of-mass energy of 1.96 TeV. The data were collected with the CDF II detector and correspond to an integrated luminosity of 5.3 fb$^{-1}$. No significant excess above the standard model expectation is observed and we set upper limits on $\sigma \cdot B(W' \rightarrow e\nu)$. Assuming standard model couplings to fermions and the neutrino from the $W'$ boson decay to be light, we exclude a $W'$ boson with mass less than 1.12 TeV/$c^2$ at the 95% confidence level.

PACS numbers: 12.60.Cn, 13.85.Rm, 14.70.Pw

The $W'$ boson is a postulated charged heavy vector boson...
son which is predicted in models that extend the gauge structure of the standard model. In the left-right (LR) symmetric model considered here, the right-handed $W'$ boson mass is obtained by the symmetry breaking of the right-handed electroweak gauge group of $SU(2)_R \times SU(2)_L \times U(1)_{B-L}$. This provides a natural explanation for the observed suppression of $V + A$ currents in low energy weak processes. The LR symmetric model can also be motivated by the manifestation of a higher symmetry predicted at intermediate energies in grand unified theories.

The manifest LR symmetric model assumes that the right-handed Cabibbo-Kobayashi-Maskawa matrix and the gauge coupling constants are identical to those of the standard model. The $W'$ can decay in the same way as the standard model $W$, with the exception that the $tb$ decay channel is accessible if the $W'$ is heavy enough and that the diboson decay channel ($W' \to WZ$) is suppressed in the extended gauge model.

The $W'$ boson has been previously searched for in high energy physics experiments using final state signatures such as leptons, jets, and/or missing energy. The most recent direct searches for a charged heavy vector boson have been performed at the Tevatron collider at Fermilab. The CDF experiment previously set limits on the cross section times branching fraction in the decay mode $W' \to tb$ and excluded a $W'$ boson mass below 800 GeV/$c^2$ at the 95% confidence level (CL) using 1.9 fb$^{-1}$ of data. The D0 experiment set limits on the product of the cross section and branching fraction in the decay mode $W' \to ev$ and excluded a $W'$ boson mass below 1.00 TeV/$c^2$ at the 95% CL using 1.0 fb$^{-1}$ of data. Both of these recent mass limits assume that the couplings between the new vector boson and the fermionic final states are the same as in the standard model.

In this Letter, we present the results of a search for a $W'$ boson in the $ev$ decay mode, assuming the manifest LR symmetric model and the right-handed neutrino from the boson decay to be light ($m_\nu \ll m_{W'}$) and stable. Under these assumptions, the results in this Letter can be useful in the generic model since the kinematics of the left- and right-handed $W'$ bosons is not different. We use a data sample corresponding to 5.3 fb$^{-1}$ integrated luminosity of $pp$ collisions at $\sqrt{s} = 1.96$ TeV recorded by the upgraded Collider Detector at Fermilab (CDF II). We select events that are consistent with the production of the standard model $W$ and the heavier $W'$ boson that decay to an electron and neutrino in the final state. The analysis technique applied is the same as in a previous search.

The CDF II detector is described in detail elsewhere. CDF II is a general purpose solenoidal detector which combines precision charged particle tracking with fast projective calorimetry and fine-grained muon detection. Tracking systems are contained inside a superconducting solenoid, 1.5 m in radius and 4.8 m in length, which generates a 1.4 T magnetic field parallel to the beam axis. Calorimeters and muon systems surround the solenoid and the tracking system. Electron candidates are identified by an energy deposit in the electromagnetic calorimeter with a track pointing to it. A set of charged-particle detectors surrounding the calorimeters identify muon candidates. The energy of the electron candidate is measured by the calorimeter and its direction is determined from the tracking system. The component of the neutrino momentum transverse to the beamline is inferred to be equal to the missing transverse energy $E_{T}$, which is derived from the transverse energy imbalance of all the deposited energy in the calorimeters.

The online selection requires either one electron candidate in the electromagnetic calorimeter with transverse energy $E_{T} > 18$ GeV that has a matching track with transverse momentum $p_{T} > 9$ GeV/$c$ or an electron candidate in the electromagnetic calorimeter with transverse energy $E_{T} > 70$ GeV. No restrictions on the amount of energy leakage into the hadronic calorimeter were imposed, in order to ensure high efficiency for high-$E_{T}$ electrons. We select the candidate event sample offline by requiring an isolated electron candidate with $E_{T} > 25$ GeV and the existence of an associated track with $p_{T} > 15$ GeV/$c$ that is contained in the fiducial region of the tracking system of $|\eta| < 1.0$.

Electron candidates are selected based on an $E_{T}$-dependent isolation cut in order to maximize the efficiency in the high-$E_{T}$ region. The electron shower profile is required to be consistent with that of test-beam electrons in order to match with the expected EM shower. In events with high-energy muons, the $E_{T}$ is adjusted by adding the muon momentum and removing the expected ionization energy deposition in the calorimeter. The $E_{T}$ is corrected further for $\eta$- and energy-dependent non-uniformities of the calorimeter response. In the final selection, the corrected $E_{T}$ is required to be greater than 25 GeV. Dilepton events coming from Drell-Yan, $t\bar{t}$, and diboson backgrounds are vetoed by rejecting events with a second isolated lepton, either an electron or a muon, with $p_{T} > 15$ GeV/$c$. QCD multijet events are a background to $W/W' \to ev$ when a jet is misidentified as an electron and mismasured jets lead to significant $E_{T}$. The electron candidate $E_{T}$ and the event $E_{T}$ are likely to significantly differ in magnitude in this case. In contrast, a $W/W' \to ev$ event will have an electron and neutrino emitted in opposite directions which results in the electron $E_{T}$ and $E_{T}$ being of comparable magnitude, respectively, assuming the $p_{T}$ of the boson is much smaller than its mass. Thus, in order to reduce the QCD
multijet background, we require the candidate events to satisfy $0.4 < \frac{E_T}{p_T} < 2.5$. The efficiency of this requirement is larger than $99\%$ for $W/W'$ events whereas the rejection rate is $\sim 40\%$ for QCD multijet events with $E_T > 100\,$GeV. After all selection requirements, the transverse mass of a candidate event is calculated as

$$m_T \equiv \sqrt{2E_T E_T (1 - \cos \phi_{e\nu})}, \quad (1)$$

where $\phi_{e\nu}$ is the azimuthal opening angle between the electron candidate and the $E_T$ direction.

![Image of transverse mass distributions for $W' \rightarrow e\nu$ signal events generated using PYTHIA with total background expectation.](image)

The background sources to $W' \rightarrow e\nu$ are primarily processes with an electron and missing energy in the final state. These sources of background are $W \rightarrow e\nu$, $W \rightarrow \tau\nu \rightarrow e\nu\nu\nu$, $Z/\gamma^* \rightarrow \tau\tau \rightarrow eX$, $t\bar{t}$, and diboson ($WW$, $WZ$) production. The $Z/\gamma^* \rightarrow ee$ process can also produce missing energy when one of the electrons escapes detection. The $m_T$ distributions and acceptance times efficiency of the non-multijet backgrounds are obtained using PYTHIA and a simulation of the CDF II detector. Theoretical cross section predictions are used to estimate the expected background yields \cite{14,17,18}. For the QCD multijet background estimation, a data-driven method is applied that uses the distribution of the azimuthal angle between the primary electron candidate and the vector sum of the jet energy. For the multijet case, a jet misidentified as an electron candidate will appear to recoil against the rest of the jet in the event. Therefore, a back-to-back distribution is expected in the azimuthal opening angle. The $W/W' \rightarrow e\nu$ process however does not have a strong correlation in this angle. The QCD multijet contribution is estimated by a likelihood fit to the data using the different angular shapes. The multijet $m_T$ distribution is obtained using a QCD enriched sideband sample with the isolation cut inverted. The data and the total background $m_T$ distributions are compared in Fig. 2. The contributions from $W \rightarrow e\nu$, QCD multijet, and the other backgrounds in the mass region above $m_T = 200\,$GeV/c$^2$ are listed in Table II. This comparison shows good agreement between the data and the total backgrounds.

In order to quantify the size of the potential signal contributions in the data sample, a binned maximum likelihood fit was performed on the observed $m_T$ distribution between 0 and $1500\,$GeV/c$^2$, using the background predictions and the expected $W'$ boson contribution for different mass values ranging from 500 to $1300\,$GeV/c$^2$. The fit results are shown in Table III normalized to

$$\beta \equiv \frac{\sigma \cdot B(W' \rightarrow e\nu)}{\sigma \cdot B(W' \rightarrow e\nu)_{LR}}, \quad (2)$$

where the numerator is the observed cross section times branching fraction and the denominator is that expected from the manifest LR symmetric model. The expected signal yield was normalized to the observed $W$ boson yield obtained from the fit. This removes several sources of systematic uncertainty such as the integrated luminosity, the trigger and the identification efficiencies, all of which cancel in the ratio.

Systematic uncertainties on the signal and the background rates were considered for the PDFs, the jet energy scale, the theoretical cross sections, the multijet background, the initial/final state radiation of the signal, and the energy scale of the electromagnetic calorimeter. The dominant contribution to the systematic uncertainty
the “physical region” where this ratio is greater than or observed to the expected cross section are obtained from function. The 95 % CL upper limits on the ratio of the nuisance parameters in the definition of the likelihood varies from events compared to the total backgrounds. To determine the limit on \( \beta \), we use a Bayesian approach by constructing a marginalized posterior probability distribution (\( p(\beta) \)) from the likelihood function. Sources of systematic uncertainty are included as nuisance parameters in the definition of the likelihood function. The 95 % CL upper limits on the ratio of the observed to the expected cross section are obtained from the fit. We use the resulting likelihood function only in the “physical region” where this ratio is greater than or equal to zero. The obtained upper limits are summarized in Table III and plotted in Figure 2 as a function \( m_{W'} \) together with the expected limits obtained from simulated experiments with background only. Using theoretical predictions that assume the manifest LR symmetric model, the limits on the cross section times branching fraction are converted into limits on the mass of the \( W' \) boson. The lower mass limit can be set at the mass value for which \( \beta_{95} = 1 \), where \( \int_0^{\beta_{95}} p(\beta)d\beta = 0.95 \). We take the lower bound of the theoretical cross section to obtain the mass limit. Hence, the 95 % CL is found to be \( m_{W'} > 1.12 \text{ TeV}/c^2 \).

| \( W \to e\nu \) | Events in \( m_{W'} \) bins (\( \text{GeV}/c^2 \)) |
|-----------------|-----------------------------|
| 200 - 250 | 350 - 350 | 500 - 500 | 700 - 700 | 900 - 900 |
| 711^{+59}_{-50} | 359^{+25}_{-25} | 85^{+8}_{-8} | 13^{+1}_{-1} | 1.1^{+0.1}_{-0.1} |
| Multijet | 9^{+2}_{-2} | 6^{+1}_{-1} | 2^{+2}_{-2} | 0.2^{+1.6}_{-0.2} | 0.01^{+1.0}_{-0.01} |
| Other background | 70^{+9}_{-6} | 33^{+4}_{-3} | 8^{+4}_{-1} | 1^{+1}_{-0.1} | 0.09^{+0.01}_{-0.01} |
| Total background | 790^{+61}_{-58} | 398^{+31}_{-30} | 94^{+9}_{-8} | 14^{+3}_{-1} | 1.2^{+2.2}_{-0.1} |
| Data | 784 | 426 | 88 | 18 | 1 |

TABLE II: The expected numbers of events from \( W' \to e\nu \) process, \( N_{\exp} \), assuming the manifest LR symmetric model and normalized by the observed \( W \) boson yield. We also show the observed relative rate of the \( W' \) boson production from the fit described in the text, and the 95% CL upper limit on this relative rate. The uncertainties are statistical only and do not include systematic uncertainties. The 95% upper limits include both statistical and systematic uncertainties.

| \( m_{W'} \) | \( N_{\exp} \) | \( \beta \left( = \frac{\sigma (W' \to e\nu) \cdot B (W' \to e\nu)_{LR}}{\sigma (W \to e\nu)} \right) \) |
|-----------------|-----------------------------|
| (\( \text{GeV}/c^2 \)) | (events) | Fit (\( \times 10^{-2} \)) | Upper Limit |
| 500 | 5828 | 0.08^{+0.21}_{-0.08} | 5.38 \times 10^{-3} |
| 550 | 3407 | 0.18^{+0.20}_{-0.18} | 7.16 \times 10^{-3} |
| 600 | 2037 | 0.28^{+0.28}_{-0.28} | 1.01 \times 10^{-2} |
| 650 | 1218 | 0.43^{+0.43}_{-0.43} | 1.52 \times 10^{-2} |
| 700 | 731 | 0.36^{+0.36}_{-0.36} | 1.52 \times 10^{-2} |
| 750 | 433 | 0.15^{+0.15}_{-0.15} | 2.80 \times 10^{-2} |
| 800 | 263 | 0.03^{+0.03}_{-0.03} | 3.82 \times 10^{-2} |
| 850 | 160 | 0.00^{+0.00}_{-0.00} | 5.68 \times 10^{-2} |
| 900 | 100 | 0.00^{+0.00}_{-0.00} | 8.97 \times 10^{-2} |
| 950 | 62 | 0.00^{+0.00}_{-0.00} | 1.49 \times 10^{-1} |
| 1000 | 41 | 0.00^{+0.00}_{-0.00} | 2.48 \times 10^{-1} |
| 1050 | 27 | 0.00^{+0.00}_{-0.00} | 4.36 \times 10^{-1} |
| 1100 | 19 | 0.00^{+0.00}_{-0.00} | 7.62 \times 10^{-1} |
| 1150 | 14 | 0.00^{+0.00}_{-0.00} | 7.62 \times 10^{-1} |
| 1200 | 10 | 0.00^{+0.00}_{-0.00} | 6.24 |
| 1250 | 8.1 | 0.00^{+0.00}_{-0.00} | 6.24 |
| 1300 | 6.7 | 0.00^{+0.00}_{-0.00} | 6.24 |

In summary, we have performed a search for a new heavy charged vector boson decaying to an electron-
neutrino pair with a light and stable neutrino in $p\bar{p}$ collisions at $\sqrt{s} = 1.96 \text{ TeV}$. We do not observe any statistically significant excess over the background expectations. We use a fit to the $m_T$ distribution to set upper limits on the production and decay rate of a $W'$ boson as a function of $m_{W'}$, and exclude a $W'$ boson with $m_{W'} < 1.12 \text{ TeV}/c^2$ at the 95\% CL, assuming the manifest LR symmetric model.

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S.
Department of Energy and National Science Foundation; the Italian
Istituto Nazionale di Fisica Nucleare; the Ministry of Education,
Culture, Sports, Science and Technology of Japan; the
Natural Sciences and Engineering Research Council of Canada; the
National Science Council of the Republic of China; the Swiss National Science Foundation; the
A.P. Sloan Foundation; the Bundesministerium für
Bildung und Forschung, Germany; the World Class University
Program, the National Research Foundation of Korea; the
Science and Technology Facilities Council and the
Royal Society, UK; the Institut National de Physique
Nucleaire et Physique des Particules/CNRS and Univer-
site Pierre et Marie Curie; the Russian Foundation for
Basic Research; the Ministerio de Ciencia e Innovación,
and Programa Consolidador-Ingenio 2010, Spain; the
Slovak R&D Agency; and the Academy of Finland.

[1] G. Altarelli et al., Z. Phys. C 45, 109 (1989); 47, 676 (E) (1990).
[2] J. C. Pati and A. Salam, Phys. Rev. D 10, 275 (1974); R. N. Mohapatra and J. C. Pati, Phys. Rev. D 11, 566 (1975); 11, 2558 (1975); G. Senjanovic and R. N. Mohapatra, Phys. Rev. D 12, 1502 (1975).
[3] R. N. Mohapatra, Unification and Supersymmetry, Springer, New York, 2003.
[4] M. A. Bég et al., Phys. Rev. Lett. 38, 1252 (1977); G. Senjanovic, Nucl. Phys. B 153, 334 (1979).
[5] We omit charge, anti-particle, and flavor-tagging neutrino symbols when representing the decay modes but assume proper charge conjugation throughout this report. For example, $W' \rightarrow \ell^+ \nu_e$, and its charge conjugate are denoted as $W' \rightarrow \ell^- e^\nu$.
[6] T. Aaltonen et al. (CDF Collaboration), Phys. Rev. Lett. 103, 041801 (2009).
[7] V. M. Abazov, et al., Phys. Rev. Lett 100, 031804 (2008).
[8] A. Abulencia et al. (CDF Collaboration), Phys. Rev. D 75, 091101 (2007).
[9] D. Acosta et al. (CDF Collaboration), Phys. Rev. D 71, 032001 (2004).
[10] Missing transverse energy, $E_T^\gamma$, is defined as the magnitude of $-\sum_{i} E_{\gamma_i}$, where $\gamma_i$ is a unit vector in the azimuthal plane that points from the beam line to the center of the ith calorimeter tower.
[11] We use a coordinate system where $\theta$ is the polar angle to the proton beam, $\phi$ is the azimuthal angle about the beam axis, and $\eta$ is the pseudorapidity defined as $-\ln(\tan(\theta/2))$. Energy (track momentum) measured transverse to the beam line is denoted as $E_T (p_T)$.
[12] The additional energy ($\equiv E_{T}^{0.4} - E_{T}^{el}$) must be less then $3 + 0.02 \times E_{T}^{el}$, where $E_{T}^{0.4}$ is the transverse energy in the cone of $\Delta R (\equiv \sqrt{(\Delta \phi)^2 + (\Delta \eta)^2}) < 0.4$ centered on the electron track and $E_{T}^{el}$ is the transverse energy of the electron candidate. Both energies are in units of GeV.
[13] D. Acosta et al. (CDF Collaboration), Phys. Rev. Lett. 94, 091803 (2005).
[14] T. Sjöstrand et al., Comput. Phys. Commun. 135, 238 (2001).
[15] H. L. Lai et al. (CTEQ Collaboration), Eur. Phys. J. C 12, 375 (2000).
[16] R. Hamberg, W. L. van Neerven, and T. Matsuura, Nucl. Phys. B 359, 343 (1991); 644, 403 (E) (2002).
[17] J. M. Campbell and R. K. Ellis, Phys. Rev. D 60, 113006 (1999).
[18] M. Cacciari et al., J. High Energy Phys. 0404, 068 (2004); N. Kidonakis and R. Vogt, Phys. Rev. D 68, 114014 (2003).
[19] T. Junk, Nucl. Instrum. Methods Phys. Res., Sect. A 434, 435 (1999)

FIG. 3: The 95\% CL limits on cross section times branching fraction as a function of $W'$ boson mass and the expected limits from the simulated experiments with background only. The region above the red dashed line is excluded at the 95\% CL. The cross section times branching fraction assuming the manifest LR symmetric model, $\sigma \cdot B(W' \rightarrow e\nu)$, is shown along with its uncertainty. The intercept of the cross section limit curve and the lower bound of the theoretical cross section curve yields $m_{W'} > 1.12 \text{ TeV}/c^2$ at the 95\% CL.