White light-informed optical properties improve ultrasound-guided fluorescence tomography of photoactive protoporphyrin IX

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Abstract. Subsurface fluorescence imaging is desirable for medical applications, including protoporphyrin-IX (PpIX)-based skin tumor diagnosis, surgical guidance, and dosimetry in photodynamic therapy. While tissue optical properties and heterogeneities make true subsurface fluorescence mapping an ill-posed problem, ultrasound-guided fluorescence tomography (USFT) provides regional fluorescence mapping. Here USFT is implemented with spectroscopic decoupling of fluorescence signals (auto-fluorescence, PpIX, photoproducts), and white light spectroscopy-determined bulk optical properties. Segmented US images provide a priori spatial information for fluorescence reconstruction using region-based, diffuse FT. The method was tested in simulations, tissue homogeneous and inclusion phantoms, and an injected-inclusion animal model. Reconstructed fluorescence yield was linear with PpIX concentration, including the lowest concentration used, 0.025 μg/ml. White light spectroscopy informed optical properties, which improved fluorescence reconstruction accuracy compared to the use of fixed, literature-based optical properties, reduced reconstruction error and reconstructed fluorescence standard deviation by factors of 8.9 and 2.0, respectively. Recovered contrast-to-background error was 25% and 74% for inclusion phantoms without and with a 2-mm skin-like layer, respectively. Preliminary mouse-model imaging demonstrated system feasibility for subsurface fluorescence measurement in vivo. These data suggest that this implementation of USFT is capable of regional PpIX mapping in human skin tumors during photodynamic therapy, to be used in dosimetric evaluations. © The Authors. Published by SPIE under a Creative Commons Attribution 3.0 Unported License. Distribution or reproduction of this work in whole or in part requires full attribution of the original publication, including its DOI. [DOI: 10.1117/1.JBO.18.4.046008]

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1 Introduction

The focus of this work was to develop and test an optimized system for image-guided fluorescence tomography of subsurface skin lesions, with the goal of providing region-based quantification of protoporphyrin IX (PpIX) fluorescence in vivo. Photodynamic therapy (PDT) based upon PpIX as the photosensitizer is an established, noninvasive dermatologic treatment for actinic keratosis in the US, and is commonly used ‘off label’ for cosmetic treatments, and treatment of squamous cell carcinoma (SCC) and superficial basal cell carcinoma (BCC).1–3 PpIX-based PDT involves the administration of a precursor molecule, either aminolevulinic acid (ALA) in the form of Levulan, or methyl aminolevulinate (MAL) in the form of Metvix. The induced PpIX accumulates preferentially in epithelial pathologies including SCC and BCC, with tumor-to-background contrast from 2:1 to 9:1 depending on the pathology and the preparation of ALA used.4–6 The photophysically active form of PpIX is fluorescent, and so the effective fluorescent yield is an indirect measurement of the active concentration available for a fixed light delivery.7

While PDT often yields successful treatment outcomes in superficial (<1 mm thick) BCCs and SCCs, treatment of subsurface pathologies at deeper locations has been less successful.8,9 This limit to effectiveness is potentially due to a reduced amount of ALA, and thus PpIX, at greater depths.10 Additional methods have been developed to increase PpIX induction including differentiation therapy to increase PpIX synthesis and contrast11 and the use of lipophilic ester variations of ALA12 or vertical micro-channel ablation13,14 to increase ALA penetration. In order to understand how the influence of these methods on PpIX concentrations within targeted tissue correlates with treatment outcomes, a method must be developed to measure in vivo PpIX fluorescence distributions during PDT treatment.15 PpIX is a natural fluorophore, so molecular imaging could provide the spatial distribution information necessary for noninvasive tumor detection and delimitation,16,17 as well as evaluation and real-time outcome prediction for PDT.

Toward this goal, diffuse fluorescence tomography (FT) of tissue can be used for subsurface fluorescence imaging, and can provide improved quantitative accuracy as compared to broad-beam surface imaging methods.18–20 Subsurface FT typically employs an array of fiber optic source-detector pairs to measure diffusely remitted light, and the data are used to construct depth images from the different light paths sampled. Ideally, each source-detector pair samples a slightly different light path through the tissue and numerical method solutions for the multiple-scattering paths are used to isolate the local contribution of fluorophore concentration to each measurement.20,21 However, mathematical reconstruction of a true, high-resolution fluorescence image from FT measurements is an ill-posed problem because the limited number of measured signals does not
provide a unique solution to the fluorescence distribution map for the large number of locations within the domain. Earlier work has shown this inverse problem to be inaccurate without additional information. The addition of prior structural segmentation information using MRI, X-ray computed tomography (XRCT), or ultrasound imaging (US) can lead to improved reconstruction of fluorescence using tomography. Segmentation of tissue geometries into a few regions of interest converts the ill-posed problem into an over-defined problem for which a reliable, unique solution can be found, at the cost of reducing fluorescence imaging resolution to a few regions of interest.

Of the available imaging modes to measure structural prior information, US provides fast, low-cost, nonionizing, high-resolution, spatial information imaging, and high frequency US-guided FT has been demonstrated to improve fluorescence recovery in transmission geometry (source and detector fibers facing on opposite sides of tissue) FT. However, US-guided FT still stands to be improved by three important developments that we combine in this reported work:

1. Use reflectance geometry (source and detection fibers on same surface). Reflectance geometry allows seamless application with an US transducer and is more relevant to large animal (human) imaging applications. Most current FT methods use transmission geometry, which works well for small animals but requires light transmission through the region of interest, and is thus less applicable to large animal (human) imaging.

2. Spectroscopic signal decoupling. Most current FT methods measure fluorescence signals using one or two wavelengths or band pass-integration measurements, and do not adequately account for contaminating signals from other fluorophores or tissue/system autofluorescence in the measured signals. Sensitivity could be improved with spectrally resolved techniques to account for multiple fluorescent signals, such as photo-bleaching products and tissue autofluorescence.

3. White light-spectroscopy-informed optical properties. The current FT reconstruction methods require simulation of tissue optical paths, which in turn requires knowledge of tissue optical properties or, in the absence of measured values, choosing tissue optical properties based on literature standards. However, tissue optical properties vary temporally and spatially ($\mu_s = 0.001$ to $0.1$ mm$^{-1}$ and $\mu_t = 0.2$ to $4.0$ mm$^{-1}$) (Ref. 34) and the values used can significantly affect the reconstructed fluorescence signal.

We present here a high-frequency ultrasound-guided fluorescence tomography (USFT) method and system, with significant improvements over the previously presented system that address the three USFT development objectives outlined above (reflectance imaging, spectral decoupling, and white light-spectroscopy measurement of optical properties). The USFT method introduced, based on reflectance geometry measurements, incorporates fluorescence spectral fitting to isolate the signal of interest from other fluorescent signals, including system and tissue autofluorescence, and incorporates white light spectral fitting to measure local tissue optical properties for more accurate light path simulation and reconstruction of regional fluorescence. We validate the USFT system and approach for regional quantification of PpIX fluorescence yield in optically homogeneous and heterogeneous tissue phantoms and present preliminary in vivo imaging using a tumor-simulating mouse model.

## 2 Materials and Methods

### 2.1 Experimental System Design and Construction

A diagram of the US FT imaging system used herein is shown in Fig. 1. The instrument was originally introduced by Gruber et al. but has undergone significant hardware and software improvements noted here. It is composed of a Sonix Tablet

![Fig. 1 Ultrasound-guided fluorescence tomography system schematic](https://www.spiedigitallibrary.org/journals/Journal-of-Biomedical-Optics/046008-2_April-2013_Vol-18(4).pdf)
(Ultrasonix, Richmond, British Columbia, Canada) high-frequency ultrasound system with a 40 MHz transducer (L40 8/12, Ultrasonix) coupled to a multi-channel optical tomography instrument. These modalities are integrated at the tissue interface through a custom dual transducer ‘v-block’ positioning system which facilitates precise positioning for imaging the same plane with US and FT sequentially, a process illustrated in Fig. 1(a). Image acquisition begins by using the US system to locate the region of interest, such as a subsurface tumor. Once this imaging plane is located, the positioning arm holding the ‘v-block’ is locked and an US image is captured. Then, the US transducer is slid away and replaced with the FT optical fiber array transducer via the sliding v-block design and all optical projections are acquired. The transducer v-block is designed such that both transducers image the same plane, allowing direct use of segmented US images for FT processing.

The optical component of the system consists of 4 source and 5 detection fiber channels alternating in a single line with a center-to-center spacing of 2.5 mm [Fig. 1(b)]. The excitation source is a 633 nm laser system with current and temperature regulation (Model 7404, Intense Co., North Brunswick, New Jersey). The laser output passes through a filter wheel (FW102C, Thorlabs) containing a 2 OD neutral density filter, a 650-nm short pass filter, and a beam stop, and is multiplexed to the four source channels via a 1 x 4 fiber switch (Piezosystem Jena, Hopedale, Massachusetts). The output of each channel is coupled to the target tissue by 600 μm fibers (LGOptics, Germany). The five detection fibers channel light from the target tissue through custom moveable filter blocks to five compact spectrometers (USB2000+, Ocean Optics, Dunedin, Florida) with spectral range of 540 to 1210 nm. Sources are illuminated sequentially, while detection fiber spectra are captured in parallel. Dark spectra are captured for each imaging mode using the same exposure times and light paths with the excitation source blocked by a beam stop in the filter wheel.

In addition to the fluorescence excitation source, a shuttered white light source (HL-2000-FHSA, Ocean Optics) has been incorporated into the optical system, through an additional transducer fiber, to obtain white light spectroscopy data for explicit recovery of tissue optical properties. This capability is unique in fluorescence tomography systems.

The fiber switch, filter wheel, and spectrometers are connected via a USB hub (DUB-H7, D-link) to the Sonix Tablet Windows-based computer. The detection-side moveable filters, fiber position feedback limit switches, and white light shutter are controlled or read via NI DAQs (NI USB 6099, National Instruments). Imaging, including source-detector-pair-dependent auto-exposure for each imaging mode is controlled via LabView virtual instrument software (LabView v8.4.2, National Instruments). Ultrasound image-segmentation, spectral processing, and fitting of fluorophore distribution were performed using Matlab (R2011b, Mathworks, Natick, Massachusetts) and NIRFAST.15

To acquire data for optical reconstruction, three sets of measurements were recorded [Fig. 2(a) to 2(c)]. Once the fibers were positioned, excitation intensity was measured for all 20 projections (four sources and five detectors) [Fig. 2(a)]. During this acquisition, the source intensity was attenuated at the filter wheel using a neutral density filter of OD = 2.0 (NE20B, Thorlabs, Newton, New Jersey) to prevent detector saturation. To acquire full projection fluorescence data, the neutral density filter at the filter wheel was replaced with a 650-nm short pass filter (FES0650, Thorlabs), and 655-nm long pass interference filters (655ELPF, Omega Optics) were automatically lowered into each detection channel via movable filter mounts (MFF001, Thorlabs) attached to custom filter blocks [Fig. 2(b)]. Finally, a beam stop in the filter wheel blocked the excitation laser, and white light spectra from a single source were measured at each detector position to provide five projections for tissue spectroscopy [Fig. 2(c)].

2.2 Data Processing

2.2.1 US image segmentation and mesh generation

Ultrasound images were manually segmented into 2 regions: tumor (fluorescent inclusion), and background bulk tissue [Fig. 2(d) to 2(f)] using Matlab. Because the US transducer field of view is narrower than the FT array, the mask background was extended laterally. This mask was supplied to a meshing routine described elsewhere6,37 and recently built into the NIRFAST package along with source and detector fiber locations, node density, and mask pixel dimensions as inputs, to produce a finite element mesh.

(a) Excitation Spectrum  (b) Fluorescence Spectrum  (c) White Light Spectrum

(d) US Image  (e) Segmented US Image  (f) Region Mask

![Fig. 2 Sample FT and US images.](image-url)
2.2.2 White light spectral analysis

Broadband white light spectra in the wavelength range 550 to 900 nm were analyzed using a fitting algorithm that used diffusion theory\(^3\) to relate wavelength-dependent estimates of the reduced scattering coefficient and absorption coefficient to the calibrated reflectance intensity measured at multiple source-detector separation distances. Measured spectra were calibrated using measured intensity (\(I_{\text{meas}}\)) and model estimates of reflectance intensity (\(I_{\text{calibration}}\)) from a phantom of known optical properties, in this study a phantom with 1% Intralipid. This calibration was given as:

\[
R_{\text{sample}} \rightarrow \frac{I_{\text{calibration}}}{I_{\text{meas}}} = f(\mu_s' , \mu_a , n_{\text{rel}}, r).
\]

This calibration factor transformed the collected intensity (\(I_{\text{meas}}\)) from counts/sec to reflectance (\(R_{\text{sample}}\)) units of photons/mm\(^2\), resolving the experimental data into the same units as exported by the standard diffusion theory model. This calibrated spectrum was described by a model as:

\[
R_{\text{sample}} = \frac{\text{model}}{\text{calibration}}.
\]

Here, the model estimate of reflectance intensity (\(R_{\text{sample}}\)) includes a wavelength-dependent background scattering model, expressed as \(\mu_s' = a \lambda^{-b}\), which allows estimation of a combined Mie- and Rayleigh-type scattering. The model also specified absorption, \(\mu_a\), as the summation of the contribution of individual chromophores, as \(\mu_a = \sum \mu_a \). The chromophores within phantoms in this study included oxygenated and deoxygenated hemoglobin and water, and analysis of in vivo data also included collagen. The influence of variation in the index of refraction (\(n_{\text{rel}}\)) between calibration phantom and sample was estimated by incorporating an amplitude factor, as in Refs. 39 and 40. The fitting algorithm in this investigation estimated scattering amplitude and slope and chromophore concentrations in order to minimize the difference between the measured and model estimated reflectance spectra at each source-detector distance measured [Fig. 3(c)]. White light spectra-based bulk estimates of \(\mu_s'\) and \(\mu_a\) at excitation (633 nm) and emission (700 nm) wavelengths were used as inputs to FT reconstructions to account for variations in bulk tissue optical properties.

2.2.3 Data processing and fluorescence reconstruction

After subtraction of dark signal, spectra were normalized with respect to exposure times. A three-point median filter was used to eliminate spikes from each spectrum followed by a 14-point (5 nm) moving average smoothing filter. A single excitation value was obtained from the 633 nm peak for each projection [Fig. 3(d)]. For fluorescence spectra, a linear-least-squares fitting algorithm\(^3\) was used to determine the contributions from normalized basis spectra for PpIX fluorescence, photoproduction photoprotoporphyrin (pPp), and system and tissue auto-fluorescence [Fig. 3(a)]. The pPp fluorescence basis spectra was approximated using a Gaussian curve centered at 675 nm with 27 nm full-width-half-maximum, as reported previously by Kruijt et al.\(^4\) The combined system and target auto-fluorescence spectra were obtained for each detection channel using the same phantom ingredients without PpIX [Fig. 3(a)]. This analysis yielded a fitted coefficient representing the fluorescence...
attributable to PpIX within the measured spectrum at each source-detector pair [Fig. 3(b) and 3(e)] and a corresponding transmitted excitation signal value [Fig. 3(d)].

The diffusion-based forward model as described by Davis et al.\textsuperscript{32,43} was used to simulate the reference surface excitation ($\Phi_{\text{ref},x}$) and expected fluorescence signals ($\Phi_{\text{calc},x}$), where subscripts $x$ and $m$ denote quantities at the excitation and emission wavelength, respectively, using target optical properties recovered with white light spectroscopy.

The measured, excitation-normalized data $\Phi_{\text{BR}}$ [$\Phi_{\text{BR}} = \Phi_{m}/\Phi_{x}$; Fig. 3(f)] was scaled by the simulated reference data ($\Phi_{\text{ref},x}$) to provide a calibrated fluorescence measurement for each source-detector pair for reconstruction with the NIRFAST architecture.

$$\Phi_{\text{fl,calib}} = \Phi_{\text{BR}} \Phi_{\text{ref},x}.$$  

The fluorescence yield was calculated using an iterative error minimization algorithm with the objective function

$$X^2 = ||\Phi_{\text{fl,calib}} - \Phi_{\text{fl,calc}}||^2,$$

where $\Phi_{\text{fl,calc}}$ is the simulated fluorescence signal [Fig. 3(f)]. The corresponding iterative update equation applied was

$$\delta \eta_{\text{diff}} = (J^T J + \lambda I)^{-1}J^T \delta \Phi,$$

where $\eta_{\text{diff}}$, the effective fluorescence yield, is the product of the fluorophore quantum efficiency $\eta$ and absorption coefficient $\mu_{\text{df}}$, $\delta \mu_{\text{df}}$ is the update parameter, $\lambda$ is the fixed regularization parameter, $\delta \Phi$ is the error between $\Phi_{\text{fl,calib}}$ and $\Phi_{\text{fl,calc}}$ at each iteration, and $J$ is the Jacobian (sensitivity) matrix.

Spatial segmentation information from the co-registered ultrasound image was incorporated into the reconstruction algorithm, thus drastically reducing the dimensionality of the problem from fitting fluorescence yield at thousands of mesh nodes to fitting fluorescence yield at just a few homogeneous regions. An extensive review of the benefits of spatial prior in FT is provided in Ref. 31.

The 2-region mesh obtained from the segmented US image was used to generate a region-mapping matrix $k$ that maps the Jacobian to a lower dimension: \textsuperscript{44,45}

$$J' \rightarrow Jk, \quad \text{where } k_{ij} = \begin{cases} 1, & \text{node } i \in \text{region } j, \\ 0, & \text{otherwise} \end{cases}$$

And the new update equation:

$$\delta \eta_{\text{diff}} = (J'^T J' + \lambda I)^{-1}J'^T \delta \Phi.$$  

Since FT reconstructions are generally capable only of recovering the effective fluorescence yield, $\eta_{\text{diff}}$, which is dependent on the quantum yield, molar extinction coefficient, and concentration of the fluorophore, actual fluorophore concentration cannot be calculated without estimates of the other properties. The quantum yield for PpIX in DMSO was determined to be $\eta = 0.0046$ using the comparative method of Williams et al.\textsuperscript{45} with Kiton Red in DMSO as a standard, but it is well known that this value varies depending on solvent and local environment. The PpIX molar extinction coefficient at 633 nm is 4866 M$^{-1}$cm$^{-1}$,\textsuperscript{46} as found on the Oregon Medical Laser Center website (http://omlc.org/). Using this molar extinction coefficient and the known PpIX concentrations, effective quantum yield was calculated for each experimental series.

### 2.3 Experimental Validation Studies

#### 2.3.1 Phantoms

In order to examine the sensitivity of the USFT system to changing PpIX concentration in a homogeneous liquid medium, the system was tested using homogeneous liquid phantoms containing Intralipid (20%, Fresenius-Kabi, Bad Homburg, Germany), 5% Tween20 (P1379, Sigma-Aldrich), porcine whole blood (7204901, Lampire Biological Inc.), and [0.025,0.05,0.1,0.2,0.4,0.8] μg/ml PpIX. Blood and Intralipid concentrations were varied to generate a sampling cross using: values for (blood%, Intralipid%) were (1%, 0.5%) (1%, 1%) (1%, 1.5%) (0.5%, 1%) (1.5%, 1%). Each phantom was stirred continuously until imaging. \textsuperscript{29,47} Three images were obtained for each combination of blood, Intralipid, and PpIX concentration, resulting in 15 measurements for each PpIX concentration, with three measurements for each of the five blood/Intralipid values. Reconstruction was performed in two stages. First, all phantoms were assumed to have equal optical properties based on values found in Ref. 34. Second, reconstructions were performed with phantom-specific optical properties, as measured by white light spectroscopy. Simulations were used to expand the range of optical properties, beyond those found in phantoms with $1 \pm 0.5$% blood (fully oxygenated), $1 \pm 0.5$% Intralipid. These simulations considered homogeneous phantoms with base optical properties of ($\mu_\text{a}, \mu_\text{sc}, \mu_\text{abs}$) = ((2.0, 0.15) or (0.02, 0.015))mm$^{-1}$, and ($\mu_\text{a}', \mu_\text{sc}', m$) = [(2.5, 2.0)mm$^{-1}$], with each optical property varied by factors of [0.5, 0.75, 1, 1.25, 1.5], and across 7 PpIX concentrations, for a total of 350 simulations. The simulated data, $\Phi_{\text{fl,calc}} + 5\%$ noise, was reconstructed using either the fixed base optical properties or the “measured” (known) optical properties with $5\%$ noise. Normalized standard deviation (std/mean), reconstruction objective function error, $E = \Sigma_{n=1}^2 \left( \log \text{data}_n - \log \text{recon}_n \right)^2$, and reconstruction linearity with concentration were calculated in each case for all optical property sets and compared to determine the effect of known optical properties on reconstructed values.

To examine the ability of the USFT system to resolve changing PpIX concentration in a small region of interest surrounded by a background with low-level fluorescence, phantoms consisting of two regions, a background and a $10 \times 40 \times 4$ mm fluorescent inclusion, were constructed using 10% w/v gelatin (G2500, Sigma-Aldrich), 0.085% w/v TiO$_2$ (14021, Sigma-Aldrich) for optical and ultrasound scattering, 1% porcine whole blood (7204901, Lampire Biological Inc.) for absorption, varying concentrations of PpIX (P8293, Sigma-Aldrich), and 5% Tween20 (P1379, Sigma-Aldrich) to prevent PpIX aggregation. Gelatin was dissolved in hot water, stirred vigorously until cooled to 30°C, then mixed with the remaining ingredients and poured into molds. Inclusions were cast using negative molds, and backgrounds were cast with positive molds positioned to leave inclusion-sized holes. The bulk background contained a constant concentration of 0.1 μg/ml (177 nM) of PpIX, while the inclusions contained [0.1,0.2,0.4,0.6,0.8,1.0] μg/ml PpIX. A 2-mm-thick skin-simulating layer having the same composition as the bulk background was prepared as well, and phantoms were imaged with and without the skin-simulating layer.
2.3.2 Animal models

For a preliminary study of application of the USFT system to measured PpIX in a region of interest in vivo, approximately 150 μl of Matrigel (#356231 BD Biosciences), with either 0 or 0.2 μg/ml of PpIX, was injected intra-dermally into the hind flank and abdomen of 6-week-old nude mice. The injections were made while drawing the syringe out, to form broad flat shapes, in an attempt to mimic BCC structure. Once solidified, injections were imaged using the USFT system with the inclusion centered in the image frame. Unpaired pre-injection images were captured and averaged for autofluorescence basis spectra. US images were segmented manually based on the hypoechoic nature of Matrigel injections and processed as detailed above. This approach to in vivo imaging was highly controllable for this validation phase of the instrument.

3 Results

A representative fluorescence spectral fit, for data measured on an inclusion phantom, is presented in Fig. 3(a). The combined target/system autofluorescence signals all follow trends similar to the autofluorescence spectra in Fig. 3(a). The excitation-light-normalized PpIX spectral fitting coefficients, termed the Born-normalized fluorescence, are shown in the bar plot in Fig. 3(b) for all nearest neighbor source-detector pairs. Fluorescence intensity from a projection that sampled the bulk of the fluorescence inclusion and a second projection that sampled just the background, as illustrated by the plotted sensitivity matrices, are highlighted. A representative white light spectral fit is shown in Fig. 3(c), with each line representing the added spectral effects of a fitted parameter. Figure 3(d) and 3(e) shows all 20 measured PpIX coefficients and all 20 measured transmitted excitation values, respectively, with the source location plotted for each five measurements (serial source illumination with parallel detection). Figure 3(f) shows the normalized (Fluorescence/Excitation) PpIX coefficients and the results reconstructed signals using the optical properties determined in Fig. 3(c). Note the reconstruction accurately matches the measured signal.

For homogeneous liquid phantoms, the bulk optical properties were calculated using white light spectroscopy and used in the reconstruction of fluorescence yield over a range of PpIX concentrations. The results were compared to effective fluorescence yield obtained using a fixed set of literature-based optical properties (μa, s = 0.2 mm⁻¹, μs, m = 0.15 mm⁻¹, μs, x = 3.0 mm⁻¹, and μa, m = 2.5 mm⁻¹) for all phantoms (Fig. 4). In Fig. 4(a) the reconstruction fitting errors (objective function) are plotted against known PpIX concentration in liquid phantoms, for reconstructions using optical properties calculated for individual phantoms using white light spectroscopy. The reconstructed fluorescence yield using measured optical properties is plotted against known PpIX concentration in Fig. 4(b). The use of the correct optical properties for the reconstruction model lead to significant reduction in reconstruction fitting errors, by a factor of 8.9 ± 5.1 [mean ± std, Fig. 4(a)], and lower normalized variation (standard deviation/mean) in reconstructed fluorescence yield by a factor of 2.0 ± 0.6 (data not shown). The reconstructed fluorescence values were linear with PpIX concentration, with R² = 0.96 for fixed literature based optical properties (data not shown) and R² = 0.98 for white light-measured optical properties [Fig. 4(b)]. Based on reconstructed fluorescence yield, PpIX molar extinction coefficient, and known concentration, the calculated quantum efficiency, η, was roughly 0.0025 for homogeneous liquid blood and Intralipid phantoms.

It is worth noting that measurement of optical properties using white light spectroscopy revealed that the optical absorption values of 1% blood, 1% Intralipid phantoms, when mixed in an open (oxygenated) environment, were surprisingly low (near the lower limit of the expected range of absorption for tissues) at the wavelengths of interest (633 and 700 nm) due to the much lower absorption of oxygenated hemoglobin in these regions.48 Thus, the homogeneous phantom series did not cover the full range of interest of optical absorption values. To investigate the importance of accurate optical properties for fluorescence reconstruction over the full range of expected tissue optical properties, signals were simulated over a range of optical properties centered at absorptions one and two orders of magnitude higher than the phantom absorption and analyzed. The normalized standard deviations and reconstruction fitting errors were significantly improved as for the imaged homogeneous phantom series. Even with 5% noise added to both the simulated fluorescence signal and the “measured” (known) optical properties, reconstruction with informed optical properties provides significant improvements over standard reconstructions in simulated homogeneous phantoms, similar to the improvements seen in the real homogeneous phantoms (Fig. 4).

Fig. 4. Recovery of PpIX fluorescence yield in homogenous liquid phantoms. (a) Comparison of reconstruction fitting errors (objective function: E = Σ (log data, − log recon)^2) between literature-based, fixed optical properties (+), and measurement-specific optical properties determined with white light spectroscopy (*). The errors are significantly lower using white light properties at all PpIX concentrations (*, p < 0.005). (b) Reconstruction performed using measurement-specific optical properties determined with white light spectroscopy. Results are linear, R² = 0.98.
Figure 5 shows a representative inclusion reconstruction and the benefits of prior spatial information and accurate optical property information. Figure 5(b) shows hot spots in the uninformed (no US guidance, ill-posed) reconstruction and the obviously poor fluorophore mapping when using only FT. Figure 5(c) compares reconstruction fits using literature-based or white light spectroscopy measured optical properties for a single inclusion. Optical properties used from the literature were based on known absorption and scattering of blood, Intralipid, and gelatin. Note the reconstruction fitting algorithm cannot adequately scale the intensity range, based largely on absorption, to reproduce measured data by changing only the fluorophore concentration. In Fig. 5(d) note that in two-region reconstruction, there are only two fluorophore values calculated when using prior spatial information: homogeneous values in the inclusion region and the background region.

Figure 6 Reconstruction of inclusion phantom fluorescence yield. Blue data is from phantoms with a 2-mm skin-simulating top layer, red data is from phantoms without a skin layer. (a) Recovered PpIX yield in the inclusion (open markers) and background (solid markers). Both reconstructions are linear, with $R^2 = 0.98$. Sensitivity, based on intercept, is $\sim 0.1 \mu g/ml$ PpIX. (b) Inclusion contrast-to-background ratio in both types of phantoms. The linear fit equations are $y = 1.8x - 0.5$, $R^2 = 0.76$, and $y = 1.2x - 1.0$, $R^2 = 0.78$, for phantoms with and without the skin-simulating top layer, respectively.
Figure 6(a) shows reconstructed fluorescence yield versus known fluorophore concentration for inclusion phantoms with and without a 2-mm skin-simulating top layer (same fluorophore concentration as background, 0.1 \( \mu g/ml \)), as well as the reconstructed background values (same for all inclusions, 0.1 \( \mu g/ml \)) for each inclusion fluorophore concentration. These reconstructions were performed using spatial priors from segmented US images [Figure 5(a)] and the actual optical properties were determined with white light spectroscopy. Reconstructed values are linear with concentration for inclusions with and without skin, with \( R^2 = 0.95 \) for both series, down to the background concentration and lowest concentration used, 0.1 \( \mu g/ml \). Figure 6(b) shows the calculated versus known contrast to background ratio (CBR), for inclusions with and without skin-simulating top layers. The linear fit equations are \( y = 1.21x - 0.95 \) (\( R^2 = 0.78 \)) and \( y = 1.78x - 0.52 \) (\( R^2 = 0.77 \)) for inclusions without and with the skin-simulating top layers, respectively. The mean error in recovered CBR is 23% without a top layer and 74% with a top layer. The CBR slope is near unity for the plain inclusion phantom, but is considerably higher when a skin-simulating layer is added. When all FT fiber pairs sample a uniform layer and spatial information is not available, the solution is nonunique because fluorophore concentration and depth are no longer independent parameters. The addition of US-guidance provides more accurate reconstruction through layers, as Fig. 6 shows, but does not completely solve this reconstruction issue. Based on reconstructed fluorescence yield, PpIX molar extinction coefficient, and known concentration, the quantum efficiency, \( \eta \), was roughly 0.0024 for gelatin-based blood and Intralipid phantoms without a skin-simulating top layer.

Finally, Fig. 7 shows a combined USFT reconstruction image for a 150 \( \mu l \) subcutaneous hind-flank injection of 0.2 \( \mu g/ml \) PpIX-matrigel solution in a nude mouse. The reconstructed fluorescence yield, combined with the molar extinction coefficient, indicated that the quantum efficiency, \( \eta \), was roughly 0.001, indicating the quantum yield calibration value may be significantly different in vivo compared to Intralipid- and gelatin-based phantoms. Contrast to background is not an informative metric for this case, as it is influenced by the near 0 background.

**Fig. 7** Preliminary in vivo data. Sample US and fluorophore reconstruction images for a 0.2 \( \mu g/ml \) 150 \( \mu l \) subcutaneous inclusion in a nude mouse. Contrast to background is less meaningful for an injected fluorophore, where the background is near 0.

### 4 Discussion

This paper presents the validation of a combined high frequency US-guided reflectance-mode-FT system, with spectral fitting to decouple multiple signals and novel use of white light spectroscopy to inform optical models for accurate reconstruction of regional, subsurface fluorescence yield in vivo. The data show that the US-guided spatial priors are critical for accurate regional fluorescence reconstruction. The data also underscore the benefit of white light spectroscopy and the necessity for knowledge of optical properties for accurate fluorescence tomographic reconstruction.

The perturbative influence that variations in background optical properties have on the reconstructions of raw fluorescence spectra is shown in Fig. 4(a). While it is possible to estimate a linear metric of fluorescence from reconstructions with assumed background optical property values, the overall reconstruction accuracy of fluorescence estimates are improved significantly when sample-specific optical properties are measured by white light spectroscopy and used to inform reconstruction models. This improvement in accuracy results in reduction in the normalized standard deviation of reconstructed fluorescence yield by a factor of \( \sim 2 \) across varying phantom optical properties and reduction in reconstruction fitting errors by a factor of \( \sim 9 \). In many imaging applications, the region of interest may be pathologically altered and thus tissue optical properties may differ from values expected in normal tissue. The data in this study suggests that incorporation of white light spectroscopy to determine these bulk properties will significantly reduce fluorescence reconstruction errors in vivo.

FT approaches that focus on diffuse recovery exclusively are often limited by an ill-posed problem of estimating both magnitude and spatial location of fluorescence origin within the sampled medium. The benefit of using US images to provide prior spatial information to overcome this limitation is made obvious in Fig. 5(b). The non-region-segmented problem is highly under-determined and the reconstruction algorithm cannot yield accurate spatial maps of subsurface fluorescence.\(^{9,20}\) Standard reconstruction typically yields sparsely distributed hot spots surrounded by relatively low values [Fig. 5(b)]. While such a result does very accurately fit the input measured data, the solution is not unique and does not provide an informative description of subsurface fluorescence distribution. When the reconstruction is constrained to a few homogeneous regions, the problem is no longer ill-posed and yields a more-accurate regional fluorescence yield, with the trade-off that spatial resolution is reduced to a few homogeneous regions.

The optical phantoms utilized in this study were constructed to mimic the clinically relevant problem of assessing PpIX fluorescence within a subsurface structure, such as a BCC or SCC. This situation has multiple complicating factors, including variable target size and depth, variable thickness of overlaying (nontargeted) tissue, and variations in the contrast ratio between the target and background. The data in this study show that for a 10-mm-wide \( \times \) 4-mm-thick fluorescent inclusion with CBR ranging from \( \sim 1.5 \) to 10, the US-guided FT method produces reasonable results, with measured fluorescence yield, \( \eta \mu_{af} \), roughly equal in the background and inclusion at 1:1 known CBR and then linearly increasing with inclusion concentration [Fig. 6(a) and 6(b)]. The addition of a skin-simulating top layer leads to an apparent overestimation of the \( \eta \mu_{af} \), though this is likely due, at least in part, to the compression of the top layer during imaging. Optical signal intensities decrease with
increasing depth, so if the 2-mm top layer is actually compressed to 1.5 mm, then the measured sublayer fluorescence signals would be higher than expected for a 2-mm layer, and the resulting reconstructions would yield an increased fluorescence yield within the inclusion and higher inclusion:background contrast.

This study presents data obtained in controlled situations using tissue phantoms and a well-behaved preclinical model. The phantoms used in this study are simple geometric models that we are able to easily segment into two or three clearly defined, uniform regions. In vivo measurements, however, will be complicated by heterogeneity in both structural (see Fig. 7) and optical maps, which will tend to decrease local fluorophore reconstruction accuracy. However, for the designed use, PDT assessment, regional fluorescence contrast is a useful metric for deciding when to treat and for evaluating treatment efficacy.

In future work, the high-resolution ultrasound images will be used to segment the imaged medium into multiple regions if necessary. Additionally, while the current white light spectroscopy routine provides estimates of bulk tissue optical properties, and does not account for heterogeneities, the workflow will be adjusted in future versions of the system to obtain estimates of optical properties specific to each source-detector measurement.

PpIX concentration is difficult to calculate because the quantum yield varies for different solvents and conditions, and cannot be accurately measured in scattering media such as tissue phantoms. So the measured fluorescence yield only provides a rough estimate of concentration. It is not possible to know the effective quantum yield in vivo, and so the results here most importantly demonstrate that USFT fluorescence yield reconstruction is linear with known concentration, and that the reconstructed contrast is relatively accurate. For the desired applications in PDT assessment and guidance, the fluorescence yield is actually a more important metric than the fluorophore concentration or quantum yield, because the treatment dose is directly related to the fluorescence yield.

Finally, an additional limitation is that reconstruction was performed using diffusion-theory-based finite element models to account for the influence of background tissue optical properties on the collected fluorescence intensity. A more accurate method would likely be to utilize the optical properties gained from white light spectroscopy to estimate the intrinsic fluorescence, and then reconstruct the fluorescence yield based on the intrinsic fluorescence instead of the collected fluorescence using either NIRFAST or a Monte Carlo-based reconstruction algorithm.

Additional in vivo tests using injected fluorophore to determine absolute fluorescence yield, concentration, contrast, region size, and region depth sensitivities will complete validation of the USFT method and system for spatial quantification of PpIX in vivo. USFT imaging will then be employed to measure the spatial production of PpIX for diagnostic purposes and for evaluation of treatment efficacy during PDT for deep skin tumors. The USFT methods presented here can also be readily applied to measure regional fluorescence of other fluorophores to quantify drug delivery, targeted binding, or molecular function.

5 Conclusion

This study provides a new hybrid imaging/optics approach to region-based molecular imaging of subsurface targets in vivo. The presented data demonstrates that ultrasound-guided fluorescence tomography (USFT) can be employed to determine region-based PpIX fluorescence at greater depths than standard surface-imaging techniques, and could provide PpIX imaging for tumor detection and delineation as well as for PDT guidance and dosimetry. The method presented here combines tissue spatial information (US) with an array of diffuse optical samples to reconstruct fluorescence yield in tumor and background regions. Fluorescence spectral fitting is used to isolate the fluorophore of interest from autofluorescence and white light spectral fitting is used to accurately inform the optical models for reconstruction. We present preliminary in vivo results supporting the claim that USFT dual imaging offers rapid, noninvasive, ionizing-radiation-free measurement of relative region-based fluorescence of a drug of interest using a portable, low-cost imaging system.

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