Supernova Magnitude Evolution and PAge Approximation

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Abstract

The evidence of environmental dependence of SN Ia luminosity has inspired recent discussion about whether the late-universe cosmic acceleration is still supported by supernova data. We adopt the ΔHR/Δage parameter, which describes the dependence of supernova absolute magnitude on the age of supernova progenitor, as an additional nuisance parameter. Using the Pantheon supernova data, a lower bound $\gtrsim 12$ Gyr on the cosmic age, and a Gaussian prior $H_0 = 70 \pm 2$ km s$^{-1}$ Mpc$^{-1}$ on the Hubble constant, we reconstruct the cosmic expansion history. Within the flat $\Lambda$ cold dark matter framework, we still find a 5.6σ detection of cosmic acceleration. This is because a matter-dominated decelerating universe would be too young to accommodate observed old stars with age $\gtrsim 12$ Gyr. A decelerating but non-flat universe is marginally consistent with the data, but, however, only in the presence of a negative spatial curvature $\sim 2$ orders of magnitude beyond the current constraint from cosmic microwave background data. Finally, we propose a more general parameterization based on the cosmic age (PAge), which is not directly tied to the dark energy concept and hence is ideal for a null test of the cosmic acceleration. We find that, for a magnitude evolution rate $\Delta HR/\Delta age \lesssim 0.3$ mag/5.3 Gyr, a spatially flat and decelerating PAge universe is fully consistent with the supernova data and the cosmic age bound, and has no tension with the geometric constraint from the observed cosmic microwave background acoustic angular scales.

Unified Astronomy Thesaurus concepts: Cosmology (343); Observational cosmology (1146); Dark energy (351); Cosmological constant experiments (335)

1. Introduction

The accelerated expansion of the late universe, one of the greatest puzzles of modern physics, was first indicated by the “unexpected extra dimming” of high-redshift SNe Ia (Riess et al. 1998; Schmidt et al. 1998; Perlmutter et al. 1999; Scolnic et al. 2018). It is explained in modern cosmology by a hypothetical dark energy component, whose microscopic nature is often interpreted as a cosmological constant $\Lambda$ (ΛCDM model where CDM stands for cold dark matter), or an unknown fluid component with a negative equation of state $w$ ($w$CDM model). Within the ΛCDM or $w$CDM framework, the late-time cosmic acceleration is also confirmed by a few independent cosmological probes, such as the cosmic microwave background (CMB; Aghanim et al. 2018) and the baryon acoustic oscillations (BAOs; Gil-Marín et al. 2016). However, CMB and BAO constraints are more model dependent, as the prediction of observables depends not only on the expansion history of the universe, but also on the growth of inhomogeneities that depends on more detailed properties of cosmic ingredients.

If the luminosity of SNe Ia can be calibrated to be a constant (with small and unbiased scattering), the Hubble diagram of supernovae would be the currently most direct and model-independent evidence for cosmic acceleration. The empirical standardization of supernova peak luminosity is obtained by a calibration against the light-curve shape and the color. The universality of such a standardization procedure is based on the assumption that SN Ia explosion is triggered by a critical condition that has little to do with its galactic environment and the past history of its progenitor. This assumption has been intensively investigated in the past decade. A series of work has found correlation between the standardized absolute magnitude of SNe Ia and properties of their host galaxies (Hicken et al. 2009; Sullivan et al. 2010; Rigault et al. 2013, 2015, 2018; Roman et al. 2018; Kim et al. 2019). In particular, Kim et al. (2019) claimed a detection of a correlation between standardized supernova luminosity and stellar population age at a 99.5% confidence. In a subsequent work, the authors interpret the correlation as a $\sim 0.27$ mag/5.3 Gyr dependence of supernova standardized magnitude on the age of its progenitor (Kang et al. 2020). Because the typical age of a supernova progenitor decreases with redshift, such a dependence dims supernovae at high redshift, and therefore is observationally degenerate with the believed accelerating expansion of the universe. For a $\sim 0.27$ mag/5.3 Gyr supernova magnitude evolution, Kang et al. (2020) showed a concrete example that the supernova Hubble diagram can be roughly fit by an open CDM universe without late-time acceleration. However, the spatial curvature parameter $\Omega_k = 0.73$ used in the example is $\sim 2$ orders of magnitude beyond the CMB constraint (Aghanim et al. 2018). Although the CMB constraint on the flatness of universe is model-dependent, a $\sim 2$ orders of magnitude boost, if ever possible, may require very careful (and fine-tuned) construction of the model.

It might be puzzling why a large $\Omega_k$ is needed to fit the Hubble diagram if supernova magnitude evolution already (at least qualitatively) mimics the effect of $\Lambda$. The real problem for a decelerating universe without $\Omega_k$ is not the detailed quantitative difference in the Hubble diagrams, but the cosmic age! A successful cosmological model must predict a cosmic age $t_0 > t_s$, where $t_s$ is the maximum age for the oldest stars. The currently most accurate astrophysical determination of $t_s$ is based on the separation of isochrones of different ages on the H-R diagram around the turn-off and subgiant branch. The recently improved parallaxes and spectra by the Hubble Space Telescope and the Gaia spacecraft give an estimation of $t_s \gtrsim 12$ Gyr, with the uncertainty reduced to a sub-Gyr level (VandenBerg et al. 2014; Catelan 2018; Sahlholdt et al. 2019). Thus, a flat CDM universe, which predicts a cosmic age $\sim 9$ Gyr, does not pass the astrophysical tests.
In summary, the supernova magnitude evolution, if confirmed by further observations, will have a significant impact on low-redshift cosmology, but it is yet unclear whether a non-accelerating universe can be made consistent with the observational facts. In this work, we will extend the qualitative discussion in Kang et al. (2020) to a full quantitative Bayesian exploration of the cosmological implication of supernova magnitude evolution. We will start with the non-flat ΛCDM model, and proceed to a more general framework beyond the usual concept of dark energy.

Throughout this Letter we use natural units $c = \hbar = 1$. A dot denotes the derivative with respect to the cosmological time $t$. The scale factor $a$ of the Friedmann–Robertson–Walker metric is normalized to unity.

### 2. Non-flat ΛCDM Revisited

We use the Pantheon supernova catalog (Scolnic et al. 2018) and modify its likelihood by adding a progenitor age modulated magnitude

$$
\Delta m = \frac{\Delta HR}{\Delta \text{age}} \times \tau_{\text{median}}(z), \tag{1}
$$

where $\tau_{\text{median}}(z)$ is the median value of the progenitor age $\tau$ for an observed supernova at redshift $z$. We assume the following priors: $0 \leq \Delta HR/\Delta \text{age} < 0.3$ mag/5.3 Gyr, $H_0 = 70 \pm 2$ km s$^{-1}$Mpc$^{-1}$ (Gaussian), and the age of the universe $t_0 > 12$ Gyr. The median age of the supernova progenitor is computed with the following probability density of finding a supernova at redshift $z$ with progenitor age $\tau$:

$$
P(\tau; z) d\tau \propto \begin{cases} 
\tau^\alpha d\tau, & \text{if } t(z) > \tau, \\
0, & \text{otherwise}
\end{cases}
$$

$$
\propto \left( t_p^\alpha + \tau^\alpha - t_p^\alpha \right) \left[ 10^{A(z'-z_0)} + 10^{B(z'-z_0)} \right], \tag{2}
$$

where $t_p = 0.2$ Gyr, $z_0 = 1.243$, $A = -0.997$, $B = 0.241$, and $z'$ is the redshift that satisfies $t(z') + t(z') = \tau$. More about the details of the recipe of supernova progenitor age can be found in Kang et al. (2020) and references therein.

The purpose of the simple exercise done here is to compare with Kang et al. (2020) from a theoretical perspective. Equation (2) should not be overly interpreted as a thorough and accurate study on the Pantheon samples. Calibration to the actual ages of the stellar systems that make up the Pantheon supernova sample, which ideally should be done, is nontrivial and beyond the scope of this tentative exploration. Noticeably, while this work is under review, a new analysis of the supernova samples used in Kang et al. (2020) claimed no evidence for supernova luminosity evolution after removal of a single poorly sampled supernova (Rose et al. 2020). Rose et al. (2020) did not find any significant residual host-age dependence for Pantheon samples after standardization.

We first perform Monte Carlo Markov Chain analysis for the flat and non-flat ΛCDM models, respectively. The results are summarized in Figure 1. For the flat ΛCDM model, the posterior of the deceleration parameter $q_0 = -0.45 \pm 0.08$ gives a $\sim 5.6\sigma$ detection of cosmic acceleration. For the non-flat ΛCDM model, a decelerating universe is marginally consistent with the data, however, at the price of introducing an enormously large $\Omega_k \gtrsim 0.5$, which is strongly disfavored by the CMB data (Aghanim et al. 2018).

The example in Kang et al. (2020), $(\Omega_m, \Omega_\Lambda) = (0.27, 0)$, is well outside the $2\sigma$ contour. In fact, the entire $\Omega_\Lambda = 0$ line is disfavored by the data, mainly because we have used a cosmological age prior $t_0 > 12$ Gyr, which was not considered in Kang et al. (2020).

### 3. The Page Approximation

The main purpose of this work is to explore the possibility of a non-accelerating universe. A proper null test of cosmic acceleration should be done in a framework beyond the concept of dark energy. We propose a very simple, yet powerful parameterization based on the cosmic age, which we dub the parameterization based on the cosmic age (PAge) approximation. The PAge approximation contains the same number of parameters as $w$CDM, but covers a broader class of scenarios beyond the usual dark energy concept.

Since the early 2000s, a Taylor expansion of the luminosity distance $d_L$ as a function of the redshift (Visser 2004)

$$
d_L(z) \approx \frac{z}{H_0} \left[ 1 + \frac{1}{2} (1 - q) z - \frac{1}{6} (1 - q - 3q^2 + j - \Omega_k) z^2 \right], \tag{3}
$$

where the “jerk” parameter $j = \frac{\dd^3a}{dt^3} |_{z=0}$, has been widely used in the literature for model-independent explorations beyond the dark energy concept. While the Taylor expansion of $z$ is convenient and is suitable for supernova data analysis at low redshift, it may fail at $z \gtrsim 1$ that is well accessible by a modern supernova catalog. Moreover, physical conditions such as $\frac{dd}{dz} > 0$ (distance increases with redshift) and $\frac{dH}{dz} > 0$ (background energy density decreases with time) are complicated in the $q$–$j$ space. Finally, because the Taylor approximation contains no information about the high-redshift universe, the cosmic age or the distance to the last scattering surface of CMB are incomputable with Equation (3). The “jerk” parameterization by design is immune to any high-redshift criticism.
Because of the aforementioned disadvantages of the local Taylor expansion, we propose instead a global approximation of the cosmic expansion history. The PAge approximation is based on two assumptions: (i) the universe is dominated by matter at high redshift $z \gg 1$ (we ignore the radiation component and the very short period before matter domination); (ii) the product of the cosmological time $t$ and the Hubble expansion rate $H$ can be approximated as a quadratic function of $t$. It can be easily shown that the two assumptions lead to

$$\frac{H}{H_0} = 1 + \frac{2}{3} \left( 1 - \eta \frac{H_0 t}{p_{\text{age}}} \right) \left( \frac{1}{H_0 t} - \frac{1}{p_{\text{age}}} \right),$$

where $p_{\text{age}} = H_0 t_0$ is the product of Hubble constant $H_0$ and the current age of the universe $t_0$, and the phenomenological parameter $\eta$ can be regarded as a quadratic fitting parameter.

There are immediately some advantages of using the PAge approximation. For instance, both of the physical conditions $\frac{dt}{dz} > 0$ and $\frac{dt}{dz} > 0$ can be guaranteed by a simple bound $\eta < 1$, and global quantities such as the observational bounds on the cosmic age can be easily applied to PAge. More importantly, the popular models in the literature—flat or non-flat, $\Lambda$CDM or $w$CDM models—can all be approximated mapped to the PAge space by matching the age of universe $t_0$ and the current deceleration parameter $d_0 = -\frac{a}{a^2}$. A few examples are given in Table 1. As shown in the last column of Table 1, such a mapping typically yields $\lesssim 1\%$ errors in distance modulus $\mu$ at $z < 1.5$, which are negligible for current supernova data analyses. For future high-precision supernova cosmology, however, the difference between physically motivated models and PAge approximation may require more careful treatment.

The loss of $\sim 1\%$ accuracy in distance modulus is compensated by an unexplored beyond-$w$CDM parameter space, as shown by the gray background color in Figure 2. Since the mapping between PAge and an effective dark energy model is approximate and the approximation becomes worse toward the gray region, the division between $w$CDM and non-$w$CDM is only in an approximate sense. Roughly speaking, the white region can be approximated with non-interacting dark energy models, while the gray region represents more complicated models, such as an interacting dark component that exchanges energy with CDM. In general, each point on the $\eta$-PAge plane should be regarded as an approximation of many physical models that share a similar expansion history.

The marginalized constraints on $\eta$ and $p_{\text{age}}$ in Figure 2 are obtained for a flat PAge universe with supernova magnitude evolution $0 \leq \Delta H_0/\Delta \text{age} < 0.3$ mag/5.3 Gyr, the Hubble constant $H_0 = 70 \pm 2$ km s$^{-1}$ Mpc$^{-1}$, and the cosmic age $t_0 > 12$ Gyr. For the dark green and light green contours we have used an additional prior on the comoving distance to the last scattering surface (13.8 Gpc $< d_A^\text{com} < 14$ Gpc) to guarantee that the theory is roughly consistent with observed CMB acoustic angular scales. The results shown in Figure 2 suggest that a decelerating PAge universe can fit the supernova data very well without obvious tension with CMB observations.

### 4. Discussion and Conclusions

The observational hints of supernova magnitude evolution may challenge the late-time acceleration and the standard flat $\Lambda$CDM paradigm. Kang et al. (2020) proposed that a non-flat universe without dark energy may roughly fit the supernova data. We did a full Bayesian analysis in this work and showed that when a cosmic age bound is applied, (i) a non-flat $\Omega_k = 0$ universe is inconsistent with the data, mainly due to the cosmic age bound; (ii) a decelerating non-flat $\Lambda$CDM universe is marginally consistent with the data, but it requires an enormously large $\Omega_k \gtrsim 0.5$ that can hardly be made consistent with CMB observations.

The $w$CDM model is another popular extension of $\Lambda$CDM. The philosophy of $w$CDM or its extensions with time-dependent $w$ is to assume simplicity in the dark energy equation of state, which may be reasonable if the dark energy concept is accepted a priori. For a null test of the cosmic acceleration, however, we need a more general description beyond the dark energy concept.

The PAge approximation is a different philosophy. It assumes simplicity in $Ht$ rather than in dark energy $w$. PAge is a phenomenological parameterization without specifying the underlying physics that drives the late-time expansion of universe. Thus, a full calculation of CMB and BAO observables requires further model constructions. Nevertheless, we assume that the flatness and acoustic angular scale constraint from CMB will remain roughly valid. In this context we find that a decelerating PAge universe is fully consistent with the data, whereas the concordance $\Lambda$CDM shown as a red dot in Figure 2 is nothing but
a good fit on the edge of the 1σ contour. The coincidental proximity \( p_{\text{age}} \approx 1 \) in the \( \Lambda \text{CDM} \) framework has inspired some recent discussion about whether we are living in a special cosmic era (Avelino & Kirshner 2016). In the much more flexible PAge framework, the viable range of \( p_{\text{age}} \) is relaxed to \( \sim [0.86, 1.00] \) (99.7% confidence), with the lower bound mainly from the astrophysical constraint and the upper bound mainly from CMB. It would be interesting to see whether the future improved astrophysical observations will push \( p_{\text{age}} \) toward \( \approx 1 \) in the PAge framework.

The BAO standard ruler inferred from the wiggling of the galaxy power spectrum in principle also can be used to constrain the background expansion of the universe. However, redshift-space distortion (RSD) and nonlinear structures in the late universe can bias the location of BAO peaks in a model-dependent way. RSD and nonlinear corrections for models far beyond the concordance \( \Lambda \text{CDM} \) can be very nontrivial. We leave the PAge exploration of BAO, as well as of many other potential probes (Wei et al. 2017; Shajib et al. 2019; Wong et al. 2019; Zhang et al. 2019; Zheng et al. 2019), as our future work.

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**References**

Aghanim, N., et al. 2018, arXiv:1807.06209
Avelino, A., & Kirshner, R. P. 2016, ApJ, 828, 35
Catelan, M. 2018, in IAU Symp. 334, Rediscovering Our Galaxy, ed. C. Chiappini et al. (Cambridge: Cambridge Univ. Press), 11
Gil-Marín, H., et al. 2016, MNRAS, 460, 4210
Hicken, M., et al. 2009, ApJ, 700, 1097
Kang, Y., Lee, Y.-W., Kim, Y.-L., Chung, C., & Lee, C. H. 2020, ApJ, 889, 8
Kim, Y.-L., Kang, Y., & Lee, Y.-W. 2019, JKAS, 52, 181
Perlmutter, S., et al. 1999, ApJ, 517, 565
Riess, A. G., et al. 1998, AJ, 116, 1009
Rigault, M., et al. 2013, A&A, 560, A66
Rigault, M., et al. 2015, ApJ, 802, 20
Rigault, M., et al. 2018, arXiv:1806.03849
Roman, M., et al. 2018, A&A, 615, A68
Rose, B. M., Rubin, D., Cikota, A., et al. 2020, arXiv:2002.12382
Sahholt, C. L., Feltzing, S., Lindegren, L., & Church, R. P. 2019, MNRAS, 482, 895
Schmidt, B. P., Suntzeff, N. B., Phillips, M. M., et al. 1998, ApJ, 507, 46
Scolnic, D. M., et al. 2018, ApJ, 859, 101
Shajib, A. J., et al. 2019, arXiv:1910.06306
Sullivan, M., Conley, A., Howell, D. A., et al. 2010, MNRAS, 406, 782
VandenBerg, D. A., Bond, H. E., Nelan, E. P., et al. 2014, ApJ, 792, 110
Visser, M. 2004, CQGra, 21, 2603
Wei, J.-J., Melia, F., & Wu, X.-F. 2017, ApJ, 835, 270
Wong, K. C., et al. 2019, arXiv:1907.04869
Zhang, Z., Gu, G., Wang, X., et al. 2019, ApJ, 878, 137
Zheng, J., Melia, F., & Zhang, T.-J. 2019, arXiv:1901.05705