The Si$_2$H radical supported by two N-heterocyclic carbenes†

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Cyclic voltammetric studies of the hydridodisilicon(0,II) borate ([Idipp][H]Si$^0$(Idipp)!![B(ArF)$_4$]), Idipp = C[N(C$_6$H$_3$-2,6-iPr$_2$)CH]$_2$, ArF = C$_6$H$_5$-3,5-(CF$_3$)$_2$, reveal a reversible one-electron reduction at a low redox potential ($E_{1/2}$ = −2.15 V vs. Fc$^+$/Fc$^0$). Chemical reduction of 1H[B(ArF)$_4$] with KC$_8$ affords selectively the green, room-temperature stable mixed-valent disilicon(0,II) hydride Si$_2$(H)(Idipp)$_2$ (1H), in which the highly reactive Si$_2$H molecule is trapped between two N-heterocyclic carbenes (NHCs). The molecular and electronic structure of 1H was investigated by a combination of experimental and theoretical methods and reveals the presence of a π-type radical featuring a terminal bonded H atom at a flattened trigonal pyramidal coordinated Si center, that is connected via a Si–Si bond to a bent two-coordinated Si center carrying a lone pair of electrons. The unpaired electron occupies the Si–Si π* orbital leading to a formal Si–Si bond order of 1.5. Extensive delocalization of the spin density occurs via conjugation with the coplanar arranged NHC rings with the higher spin density lying on the site of the two-coordinated silicon atom.

1. Introduction

Open-shell silicon hydrides are of significant importance as transient intermediates in the chemical vapor deposition (CVD) of silicon or silicon-containing thin films, which are extensively used in the semiconductor industry.1 Fundamental species in the gas phase include the SiH$_4$ ($x$ = 1–3) and Si$_2$H$_x$ ($x$ = 1–5) molecules as well as higher aggregated Si$_n$H$_m$ clusters, which are formed from silane (SiH$_4$) or disilane (Si$_2$H$_6$) in a complex cascade of reactions.1 These species, which are also of interest in astrochemistry, are unstable under terrestrial conditions and can only be detected by spectroscopic or mass spectrometric techniques.1 One scarcely studied species in this context is the Si$_2$H molecule, which was so far only detected by vibrationally-resolved photoelectron spectroscopy of Si$_2$H$^+$ anions.2 Quantum chemical calculations of Si$_2$H suggest two almost iso-energetic, $C_{2v}$-symmetric H-bridged structures, in which the unpaired electron occupies either the Si–Si π-bonding orbital ($^3B_1$ state) or a σ-type molecular orbital corresponding to the in-phase combination of the Si lone pair orbitals ($^3A_1$ state).2

Recently, N-heterocyclic carbenes (NHCs) were found to be particularly suitable Lewis bases for the thermodynamic and kinetic stabilization of highly reactive, unsaturated, low-valent Si species, leading to the isolation of a series of novel compounds with intriguing synthetic potential.3,4 Several CAAC-stabilized open-shell silicon compounds (CAAC = cyclic alkyl(amine)carbene) were also reported, in which the unpaired electron is mainly located on the CAAC substituent.4 Trapping of Si$_2$H by NHCs appeared therefore an achievable, albeit very challenging goal, given the fact that isolable molecular hydrides of silicon in an oxidation state <2 are very rare3,5 and open-shell congeners presently unknown. In comparison, three-coordinate Si$^0$ hydrides6 and four-coordinate Si$^0$ hydrides of the general formula [LB]SiH(X)(LA) (LB = neutral Lewis base; LA = neutral Lewis acid; X = singly bonded substituent)7 are meanwhile well documented.

2. Results and discussion

The hydridodisilicon(0,II) salt [[Idipp][H]Si$^0$(Idipp)]+[B(ArF)$_4$]$_2$ (1H[B(ArF)$_4$]), Idipp = C[N(C$_6$H$_3$-2,6-iPr$_2$)CH]$_2$, ArF = C$_6$H$_5$-3,5-(CF$_3$)$_2$, which was isolated recently in our group upon protonation of Si$_2$(Idipp)$_2$ (1),8 appeared to be a suitable starting material to tackle the problem of isolating an NHC-trapped Si$_2$H radical. Quantum chemical studies revealed the same sequence of frontier orbitals in 1H$^+$ and its isobal phosphorus counterpart [R$_3$P$^+$=PR]$^+$, according to which the HOMO–1
corresponds to the lone-pair orbital at the two-coordinated E atom (E = Si, P), the HOMO is the E=E \pi-bonding orbital and the LUMO is the E=E \pi^{*} orbital.\(^a\) This isolobal interrelation-ship suggested that \(1H^{*}\) might be also reversibly reducible as the phosphinylphosphinum cation \([Mes^{*}(Me)P\equiv PMes^{*}]\) \((\text{Mes}^{*} = \text{C}_{5}\text{H}_{5}-2,4,6-\text{tBu})\).\(^{11}\) In fact, cyclic voltammetric (CV) studies of \(1H[B(\text{ArF})_{4}]\) in fluorobenzene at room temperature revealed a reversible one-electron reduction at a rather low half-wave potential \((E_{1/2})\) of \(-1.63\) V as well as an irreversible oxidation at \(+0.67\) V versus the \([\text{Fe}(\text{n}^{5}\text{C}_{3}\text{Me}_{3})_{2}]^{1/2}\) reference electrode (Fig. 1 and ESI\(^{\dagger}\)).\(^{14}\) The methyl analogue \(([\text{Idipp}](\text{Me})\text{Si}^{\dagger}\equiv \text{Si}^{\ddagger}[\text{Idipp}])\) \((1\text{Me}^{*}[B(\text{ArF})_{4}])\)\(^{9}\) was found also to undergo a reversible one-electron reduction, albeit at a more negative potential \((E_{1/2} = -1.85\) V) than \(1H[B(\text{ArF})_{4}]\). Notably, reduction of \(1H^{*}\) and \(1\text{Me}^{*}\) occurs at much lower potentials than that of the cation \([\text{Mes}^{*}(\text{Me})P\equiv P\text{Mes}^{*}]\) \((E_{1/2} = -0.48\) V).\(^{13}\) This marked difference in the redox potentials of the Si- and P-based cations can be rationalized with the large increase of the LUMO energy occurring upon replacement of the two P\text{Mes}^{*} fragments by the much less electronegative isolobal Si\text{Idipp} fragments as suggested by quantum chemical calculations.\(^9\)

The CV results prompted us to attempt a chemical one-electron reduction of \(1H[B(\text{ArF})_{4}]\). Indeed, vacuum transfer of THF to a 1 : 1 stoichiometric mixture of \(1H[B(\text{ArF})_{4}]\) and \(\text{KC}_{8}\) at \(-196\) °C followed by warming to \(-40\) °C resulted in a distinct color change of the dark red solution of \(1H[B(\text{ArF})_{4}]\) to give an intensely dark green solution, which after work-up and crystallization from \(n\)-hexane at \(-60\) °C afforded \(\text{Si}_{2}(\text{H})(\text{Idipp})_{2}\) \((1H)\) as a dark green, almost black crystalline solid in 55% yield (Scheme 1) (see ESI\(^{\dagger}\)). Compound \(1H\) is extremely air-sensitive and immediately decolorizes upon contact with air, but can be stored under an atmosphere of argon at \(-30\) °C without any color change or signs of decomposition in its EPR spectrum. Thermal decomposition of \(1H\) in a vacuum-sealed glass capillary was detected upon melting at \(147\) °C leading to a dark red mass. Analysis of the soluble part of the melting residue in \(\text{C}_{6}\text{D}_{6}\) by \(^{1}H\) NMR spectroscopy revealed the presence of Idipp (95%) and \(1\) (5%).

Notably, the redox potential of \(1H[E_{1/2}\text{in C}_{6}\text{H}_{5}\text{F} = -2.15\) V vs. \([\text{Fe}(\text{n}^{5}\text{C}_{3}\text{H}_{3})_{2}]^{1/2}/[\text{Fc}/\text{Fc}]^{15}\) lies in-between that of the benzophenone radical anion \((E_{1/2}\text{in THF} = -2.30\) V vs. \(\text{Fc}/\text{Fc})^{16}\) and \([\text{Co}(\text{n}^{5}\text{C}_{3}\text{Me}_{3})_{2}]^{1/2}\) \((E_{1/2}\text{in MeCN} = -1.91\) V vs. \(\text{Fc}/\text{Fc})^{16,17}\) indicating that the radical \(1H\) is a very strong one-electron reducing agent. Consequently, the radical \(1H\) is selectively oxidized back to \(1H[B(\text{ArF})_{4}]\) upon treatment with one equivalent of \([\text{Fe}(\text{n}^{5}\text{C}_{3}\text{Me}_{3})_{2}]\) \([\text{B}(\text{ArF})_{4}]\) in THF-\(d_{8}\) (see ESI\(^{\dagger}\)). Thereby, the redox pair \(1H/1H\) provides a very rare example of a chemically reversible Si-based redox system.\(^7,17\)

Compound \(1H\) is well soluble in \(n\)-hexane, benzene, diethyl ether or THF affording intensely dark-green solutions, even at low concentrations. The origin of this intense color was analyzed by UV-Vis-NIR spectroscopy of \(1H\) in \(n\)-hexane (Fig. 2, left and ESI\(^{\dagger}\)), which revealed electronic absorptions in the whole spectral range from 220–1100 nm. Six absorption maxima were located at 254 (9970), 305 (8140), 436 (5170), 608 (7110), 704 (6860) and 958 (1440) nm, of which the intense absorptions at 608 and 704 nm are responsible for the green color of \(1H\) (the values of the molar absorption coefficients \(\varepsilon_{3}\) are given in brackets in L mol\(^{-1}\) cm\(^{-1}\)). The UV-Vis-NIR spectrum was also analyzed by time-dependent density functional theory (TDDFT) calculations (see ESI, Fig. S21\(^{\dagger}\)).\(^{18}\)

Magnetic susceptibility measurements of solid \(1H\) from 300.0–1.9 K suggest the presence of a paramagnetic compound with one unpaired electron following Curie’s law. A plot of the reciprocal molar magnetic susceptibility \((x_{m}^{-1})\) against the absolute temperature \((T)\) showed a linear correlation from which the effective magnetic moment \(\mu_{\text{eff}}\) was calculated after linear regression and found to be 1.68 \(\mu_{B}\) (Fig. 2, right and ESI\(^{\dagger}\)). This value is slightly lower than the value derived from the spin-only formula for one unpaired electron \((\mu_{\text{eff}} = 1.73\) \(\mu_{B}\)).

The molecular structure of \(1H\) was determined by single crystal X-ray crystallography. The radical features a crystallographically imposed inversion symmetry (space group: \(P2_{1}/c\)) in marked contrast to the \(C_{2}\)-symmetric structure of \(1H^{*}\) in \(1H[B(\text{ArF})_{4}]\).\(^{9}\) The Si-bonded H atom was located in the difference Fourier map and anisotropically refined with a site occupancy of 1/2 at each Si atom. However, the exact position of this H atom could not be deduced by X-ray crystallography, since structural refinements with either a terminal (Si–H) or a bridging (Si–H–Si) position led to identical \(\omega R_{2}\) values. \(1H\) features as \(1H[B(\text{ArF})_{4}]\) and \(1\) a trans-bent planar \(\text{C}_{\text{NHCC}}\)-Si–Si–\(\text{C}_{\text{NHCC}}\) core (Fig. 3). However, distinct structural differences become apparent upon comparing the three structures. For example, the Si–Si bond
1H (2.281(3) Å) is considerably longer than that in 1H[B(ArF)4] (2.1873(8) Å) or 1 (2.229(1) Å) (Table 1), and lies in-between of a typical Si=Si double bond (2.20 Å) and a Si–Si single bond (eq. 2.352 Å in z-Si). In comparison, the Si–CNHC bonds in 1H (1.873(4) Å) are shorter than the Si–CNHC bonds of the dicoordinated Si atoms in 1H[B(ArF)4] (1.940(2) Å) and 1 (1.927(1) Å) (Table 1), and similar to that of the trigonal-planar coordinated Si atom in 1H[B(ArF)4] (1.882(2) Å). Reduction of 1H+ results also in a distinct change of the conformation of the NHC substituents. Thus, both N-heterocyclic rings in 1H are arranged coplanar with the trans-bent CNHC–Si–Si–CNHC core as evidenced by the dihedral angle θCNHC of 3.3(2)° (Table 1), whereas in 1H+ one of the two N-heterocyclic rings (bonded to the two-coordinated Si atom) adopts an almost orthogonal orientation (Table 1). All these structural changes can be rationalized by quantum theory (vide infra). Thus, reduction of 1H+ leads to a population of the Si=Si π* orbital with one electron, reducing thereby the formal Si–Si bond order from 2 in 1H+ to 1.5 in 1H as nicely reflected in the computed Si–Si Wiberg bond indexes (WBI; WBI(Si–Si) of 1H+ ¼ 1.70; WBI(Si–Si) of 1H ¼ 1.17) (see ESI, Tables S11 and S12†). The coplanar arrangement of the N-heterocyclic rings allows for an optimal in-phase interaction (π-conjugation) of the Si–Si π* orbital with π*(CN3) orbitals of the NHC substituents in the SOMO of 1H (Fig. 6), providing a rationale for the shortening of the Si–CNHC bonds and the concomitant elongation of the CNHC–NHC bonds of 1H versus 1H+ (Table 1).

IR spectroscopy proved to be a very useful method to determine unequivocally the position of the Si-bonded H atom. In fact, the ATR FT-IR spectrum of 1H displayed a r(Si=H) absorption band at 2089 cm⁻¹, which is characteristic for stretching vibrations of terminal Si–H bonds (see ESI, Fig. S4†). In comparison, the r(Si–H–Si) band of SiI-H is predicted at significantly lower wavenumbers (1592 cm⁻¹ (2A1 state); 1491 cm⁻¹ (2B1 state)), and also the r(Si–H–Si) absorption bands of H-bridged silylium ions are shifted to much lower wavenumbers (ca. 1750–1950 cm⁻¹; e.g. 1900 cm⁻¹ in [Et3Si–H–SiEt3] [CHB2Cl2]1) compared with the r(Si–H) bands of the corresponding silanes (ca. 2150 cm⁻¹).22 Notably, the r(Si–H) absorption band of 1H appears in-between that of 1H[B(ArF)4] containing a trigonal planar coordinated Si atom (r(Si–H) ¼ 2142 cm⁻¹), and the Si(µ)-hydride (IMe4)SiH(SiBu3) containing a strongly pyramidal bonded Si atom (IMe4 ¼ C[N(Me)CMe3]; r(Si–H) in KBr ¼ 1984 cm⁻¹).14c Apparently, the r(Si–H) frequency decreases with increasing pyramidalization of the Si atom, which according to the quantum chemical calculations can be traced back to the decreasing s-character of the Si hybrid orbital in the Si–H bond (see ESI, Tables S11 and S12†).

Further insight into the structure of 1H was provided by continuous wave (cw) EPR spectroscopy at X-band frequencies. Spectra with a nicely resolved hyperfine coupling pattern could be obtained from samples of 1H in n-hexane solution at 336 K (Fig. 4; see also ESI, Fig. S10† for EPR spectra at different temperatures). Notably, a similar EPR spectrum was obtained in diethyl ether solution at 298 K (see ESI, Fig. S12†), suggesting that solvent coordination effects are negligible. The EPR spectrum of 1H displays a multiplet at a giso value of 2.00562, which could be well simulated assuming coupling of the unpaired electron to one 1H (I ¼ 1/2) nucleus, two different 29Si (I ¼ 1/2) and two pairs of two magnetically equivalent 14N (I ¼ 1) nuclei, respectively (Fig. 4). These observations suggest that 1H has...

Fig. 2 Left: UV-Vis-NIR spectra of 1H in n-hexane from 220–1100 nm at different concentrations (c) and path lengths (d). Right: Plot of the reciprocal molar magnetic susceptibility (Xm⁻¹) against the absolute temperature (T) (dotted black line) and the corresponding linear (red) and line equation obtained by linear regression.

Fig. 3 DIAMOND plot of the molecular structure of 1H in the single crystal at 123(2) K. Thermal ellipsoids are set at 30% electron probability. The Si-bonded H atom was omitted due to its uncertain position. Selected bond lengths [Å], bond angles [°] and torsion angles [°]: Si–Si# 2.281(3), Si–C1 1.873(4); C1–Si–Si# 109.5(1); C1–Si–Si#–C1# 180.0(3).
Table 1  Comparison of selected bonding parameters of 1H, 1H[B(Ar){SiCNHC}]

|     | Si-Si [Å] | Si-CNHC [Å] | CNHC-NCNHC [Å] | CNHC-Si-Si [°] | φ_{CNHC} [°] |
|-----|-----------|-------------|----------------|----------------|--------------|
| 1H  | 2.281(3)  | 1.873(4)    | 1.381(4), 1.402(4) | 109.5(1)       | 3.3(2)       |
| 1H[B(Ar)]_{Si-CNHC} | 2.1873(8) | 1.882(2) [Si1-CNHC] | 1.356(2), 1.358(2) | 116.73(7) [C1-Si1-Si2] | 87.11(5) |
| 1"  | 2.229(1)  | 1.940(2) [Si2-CNHC] | 1.356(2), 1.358(2) | 95.34(6) [C28-Si2-Si1] | 71.06(6) [φ_{CNHC}] |

a Data taken from ref. 9. Connectivity: \([\text{CNHC}1]\text{(H)Si1} \equiv \text{Si2}(\text{CNHC2})]\). b Data taken from ref. 19. c φ_{CNHC} denotes the dihedral angles between the CNHC-Si-Si-CNHC least-square plane and the respective N-heterocyclic ring least-square planes.

Fig. 4  Experimental (red curve) and simulated (green curve) X-band EPR spectra of 1H in n-hexane at 336 K; the ordinate (d_{iso}) is omitted for clarity. g_{iso} = 2.00562, a^{(29)Si} = 1.725 mT, a^{(77)Si} = 0.431 mT, a^{(14)N} = 0.246 mT, a^{(29)N} = 0.100 mT, a^{(1)}H = 0.605 mT.

Quantum chemical calculations of 1H were carried out using the B3LYP functional in combination with the 6-31G** basis set for the Si, N, Si-bonded H and NHC ring C atoms and the 6-31G* basis set for the peripheral C and H atoms or the B97-D3 functional in combination with RI-JCOSX approximations and the def2-TZVP basis set for all atoms. The levels of theory are abbreviated in the following with B3LYP/I and B97-D3/II. Remarkably, calculations at the B3LYP/I level of theory yielded one minimum structure (1H_{calc}), whereas two almost degenerate minimum structures were obtained at the B97-D3/II level of theory (1H_{calc} and 1H'_{calc}) (Fig. 5). All calculated minimum structures display a terminally bonded H atom bound to the Si1 atom. No minimum structure with a bridged H atom was found on the potential energy hypersurface of 1H at both levels of theory. The geometrical parameters of the minimum structure calculated at the B3LYP/I level of theory and the global minimum structure at the B97-D3/II level of theory are almost identical (Table 2 and ESI, Table S8†). These structures (1H_{calc}) contain a trigonal-pyramidal coordinated Si1 atom with a sum of angles of 335.5° (B3LYP/I) and 342.58° (B97-D3/II), respectively. Remarkably, the calculated structure of the diphosphanyl radical P2(Me)Mes-Si, which is isolobal to 1H, displays a trigonal pyramidal geometry at the three-coordinated P atom (sum of angles: 337.5°),13 as found for 1H_{calc}. In comparison, the second minimum structure obtained at the B97-D3/II level of theory (1H'_{calc}) is only 5.5 kJ mol^{-1} higher in energy than 1H_{calc} and contains the Si1 atom in a trigonal planar environment (sum of angles: 359.61°). A comparison of the structural parameters of 1H_{calc} and 1H'_{calc} with those obtained by single crystal X-ray diffraction reveals a good agreement of the calculated Si–Si, Si–CNHC and C–CNHC bond lengths of both minimum structures (Table 2 and ESI, Table S8†). While the experimental results did not allow to clearly distinguish whether a flattened trigonal-pyramidal or a trigonal-planar geometry of the H-bound Si atom is present in 1H, the theoretical studies suggest a flat energy hypersurface for the planarization of the three-coordinated Si atom.

The calculated quasi-restricted orbitals (QROs) of 1H_{calc} at the B3LYP/I level of theory and of 1H_{calc} and 1H'_{calc} at the B97-D3/II level of theory are almost identical (Fig. 6 and ESI, Fig. S17–S19†). The SOMO is the Si=Si π* orbital, confirming
that reduction of 1H⁻ leads to a population of the empty Si=Si π* orbital of 1H with one electron (see ESI, Fig. S16†). The SOMO reveals significant contributions of π* NHC orbitals due to π-conjugation. The two lower lying doubly occupied molecular orbitals (DOMOs) are the Si=Si π and the n(Si) lone pair orbital, respectively.

Notably, CASSCF(3,3)/def2-TZVP calculations26 of 1Hcalc revealed that the overall wave function is described by a major ground state configuration of [2-1-0] of the DOMO, SOMO and LUMO with 96% contribution, suggesting that static correlation can be neglected in the electronic description of 1H (see ESI†).

The calculated spin densities of 1Hcalc and 1Hcalc at the B3LYP/D3/II level of theory are depicted in Fig. 7. Mulliken analyses27 of the spin densities reveal that the highest spin density is located at the disordered Si2 atom (37% in 1Hcalc, 29% in 1Hcalc), whereas the spin density at the Si1 atom is quite small (9% in 1Hcalc, 6% in 1Hcalc), which is in full agreement with the observation of one large and one small a(29Si) hfcc in the experimental EPR spectrum of 1H (vide supra) (see ESI, Table S9f).28 Remarkably, a significant amount of the spin density is delocalized into the CNHC and N(NHC) atoms of the Si1-bonded (17% in 1Hcalc, 27% in 1Hcalc) and Si2-bonded (29% in 1Hcalc, 30% in 1Hcalc) NHC substituents, which explains the EPR-spectroscopic detection of two a(4N) hfcc’s. The calculated giso values of 1Hcalc (2.00483) and 1Hcalc (2.00454) agree well with the experimentally obtained giso value (2.00562).

Further insight into the electronic structure of 1H was provided by a natural bond orbital (NBO) analysis at the B3LYP/I level of theory (see ESI, Table S12f).28a The Si–Si bond is composed of a Si–Si σ bond and a Si=Si π bond with an occupancy of 1.95 and 0.82 electrons, respectively, which indicates indirectly a population of the Si=Si π* orbital with one electron leading thereby to a decrease of the formal Si–Si bond order from 2 in 1H to 1.5 in 1H (vide supra). The Si2 atom in 1Hcalc bears a lone pair of high s-character (72%) as similarly found for 1Hcalc (75%). Remarkably, both Si–C(NHC) bonds in 1Hcalc are composed of one doubly occupied Si–C(NHC) σ NBO and one singly occupied Si=C(NHC) π NBO, of which the latter is absent in 1Hcalc. These additional Si–C(NHC) π contributions rationalize the shortening and strengthening of the Si–C(NHC) bonds in 1H, which is also reflected in the higher Si–C(NHC) WBI indexes (1H: WBI(Si–C(NHC)) = 1.01 and 0.95; 1H: WBI(Si–C(NHC)) = 0.86 and 0.74).

Comparative analyses of the charge by natural population analyses (NPA) of 1Hcalc and 1Hcalc at the B3LYP/I level of theory reveal that the positive partial charges at the Si atoms of 1Hcalc (q(Si1) = 0.27e, q(Si2) = 0.21e) are decreased by the reduction (1H: q(Si1) = 0.14e, q(Si2) = 0.03e) (see ESI, Table S13t). Furthermore, the one-electron reduction leads to a significant decrease of the overall charges of the NHC substituents (1Hcalc: q(NHC1) = 0.36e, q(NHC2) = 0.30e; 1H: q(NHC1) = 0.05e, q(NHC2) = −0.04e), whereas the hydridic character of the Si1-bonded H atom is retained (1Hcalc: q(H) = −0.14e; 1H: q(H) = −0.18e).

**Table 2** Comparison of selected experimental and calculated bonding parameters of 1H, 1Hcalc and 1Hcalc

|          | Si1–Si2 [Å] | Si1–C1 [Å] | Si2–C2 [Å] | ∑Si [Å] | C1–Si1–Si2–C2 [Å] | φNHC1d | φNHC2d |
|----------|-------------|-------------|-------------|---------|-------------------|--------|--------|
| 1H       | 2.281(3)    | 1.873(4)    | 1.873(4)    | —       | 180.0(3)          | 3.2(2) | 3.2(2) |
| 1Hcalc   | 2.339       | 1.885       | 1.907       | 335.51  | 173.69            | 32.71  | 1.26   |
| 1Hcalc   | 2.308       | 1.861       | 1.884       | 342.58  | 173.63            | 21.95  | 3.41   |
| 1Hcalc   | 2.289       | 1.841       | 1.886       | 359.61  | 179.32            | 6.68   | 3.24   |

*Calculated at the B3LYP/6-311G***/6-31G* level of theory. **Calculated at the B97-D3/RJ-JCOSX/def2-TZVP level of theory. † ∑Si is the sum of angles around the Si1 atom. ‡ φNHC1 and φNHC2 denote the dihedral angles between the least-square plane of the atoms C1, Si1, Si2, C2 and the least square plane of the heterocyclic ring atoms of the NHC substituent bonded to Si1 and Si2, respectively.

![Fig. 6](Image) Fig. 6 Quasi-restricted orbitals (QROs) of 1Hcalc (B97-D3/RJ-JCOSX/def2-TZVP) and their corresponding energy eigenvalues; isosurface value: 0.04 e bohr⁻³; DOMO = doubly occupied molecular orbital, SOMO = singly occupied molecular orbital, LUMO = lowest unoccupied molecular orbital.

![Fig. 7](Image) Fig. 7 Spin densities of the calculated minimum structures 1Hcalc (left) and 1Hcalc (right) at the B97-D3/RJ-JCOSX/def2-TZVP level of theory. The N-bonded 2,6-diisopropylphenyl substituents are omitted for clarity.
3. Conclusions

The isolation and full characterization of NHC-trapped Si2H (1H) can be considered as a major advance in low-valent silicon hydride chemistry, given the intermediacy of Si2H in the chemical vapor deposition of amorphous hydrogenated silicon that is widely used in solar cell and thin film technologies. Whereas Si2H features a $C_{2v}$-symmetric H-bridged ground state structure and is a σ-type radical with a symmetric distribution of the spin density over the two silicon atoms, its NHC-trapped counterpart $\text{Si}_2(\text{H})\text{(Idipp)}_2$ (1H) features a terminal Si–H bond and is a π-type radical, in which the unpaired electron occupies the $\text{Si}=\text{Si}$ π* orbital (SOMO), leading to a formal Si–Si bond order of 1.5. Significant delocalization of the spin density into the NHC substituents occurs via π-conjugation of the Si–Si π* orbital with the π* orbitals of the coplanar arranged N-heterocyclic rings leading to a stabilization of the radical, in which the spin density is higher at the NHC substituents. The mixed valent disilicon(0, I) hydride chemistry, given the intermediacy of Si2H in the CVD technology, is widely used in solar cell and thin film technology.

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