Anti-reflection cover to control acoustic intensity in *in vitro* low-intensity ultrasound stimulation of cells

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**Abstract** – Low-intensity ultrasound stimulation is a technique used in therapeutic ultrasound for bone regeneration. However, the underlying mechanisms are still poorly understood. *In vitro* studies on cell cultures are implemented to understand the processes involved. To analyze the effects of ultrasonic waves on cells, the control of the delivered acoustic intensity is essential. However, depending on the sonification protocol chosen, multiple reflections and standing waves that form inside the culture medium strongly hinder the estimates. In this work, we propose the development and the experimental validation of an anti-reflection cover. We demonstrate that this custom-designed device is effective in avoiding multiple reflections and makes it possible to artificially replace the layer of culture medium with a large amount of water. Finally, an analytical study of the acoustic intensity delivered to the cells is proposed.

**Keywords:** Ultrasound stimulation, Cell-free culture medium, Acoustic intensity, Anti-reflection cover

1 Introduction

Therapeutic ultrasound [1–3] refers to the use of mechanical acoustic waves for treating a variety of biological tissue pathologies. More precisely, these acoustic waves can be classified into three categories depending on their intensities: waves of extreme intensities, namely shock waves (>10 kW/cm²) used, for example, in lithotripsy or histotripsy; high-intensity waves (1–10 kW/cm²) used, for example, for the destruction of tumors by High-Intensity Focused Ultrasound (HIFU); and low-intensity waves (<1 kW/cm²) capable of stimulating cell growth and assisting, for example, in the healing of fractures. This work focuses on this latter application of therapeutic ultrasound.

Low-intensity ultrasound stimulation (LIUS) of bone regeneration has been demonstrated since the 1950s. There are several small and portable devices, such as the commercial Exogen™ device (Bioventus Inc, Durham, NC, USA), that patients can apply directly to the skin over the site of a bone fracture. Both fundamental research and clinical studies have been conducted to investigate the mechanical and biological effects of ultrasonic waves on cells and tissue [4], yet some of the underlying physical mechanisms involved are still poorly understood. The effectiveness of the treatment using LIUS is still discussed today [5–7]. Laboratory tests with controlled [8] and/or specific experimental setups [9, 10] are performed on cells grown in a Petri dish placed above or below a transducer. Although the transducer-cell distance and position can be adapted from one configuration to another, depending on the range of ultrasonic frequencies and the shape and orientation of the beam, reflections generally occur between the bottom of the Petri dish and the upper interface formed by the air, creating standing waves in the cell culture medium. The formation of standing waves inside the Petri dish of the *in vitro* LIUS device is a challenging problem that prevents accurate prediction of exposure conditions of the cells. This limitation of numerous *in-vitro* ultrasound stimulation protocols is often mentioned in the literature [11–13] and some technical solutions are proposed [14–17] but without examining their effectiveness, except a very recent paper [10].

2 Context & objectives

Our work focused on the development and the experimental validation of an anti-reflection cover placed over a Petri dish, under LIUS conditions in terms of frequencies,
acoustic signal signatures, pressures and intensities. The Petri dish contained a cell-free degassed water layer as a hypothetical culture medium.

The tests were conducted using immersion transducers with center frequencies of 1 MHz and 2.25 MHz that were calibrated prior to the anti-reflection cover tests. These calibrations were performed in short-pulse mode, and the cover tests were performed in tone-burst mode under LIUS conditions. For both transducers, we studied the acoustic intensities with and without the Petri dish to verify intensity distributions and levels. Reflection and transmission coefficients, peak acoustic pressures, and acoustic intensities were measured and calculated for the different configurations. The custom-designed anti-reflection cover proposed in this work reduced the reflection of the waves at the interface between water and air, trapped the transmitted waves, and thus prevented the standing waves in the Petri dish water layer. We investigated how the cover over the Petri dish can influence the peak acoustic pressure. We showed that the acoustic intensity distribution and level obtained with the cover on the Petri dish containing a millimetric water layer are equivalent to those which would be measured if the Petri dish contained a high water column of several centimeters in height. We demonstrated that it is possible to analytically estimate the energy transmitted into the Petri dish and received by the cells exposed to a low intensity ultrasound field.

3 The Petri dish/anti-reflection cover set

Figure 1 shows the details, dimensions and pictures of the Petri dish/anti-reflection cover set.

The Petri dish used in the set was a Corning model (ref. 430 196) with a diameter of 60 mm, a depth of 15 mm and a wall thickness of 0.8 mm. The material the Petri dish is made of was polystyrene, with a mass density ($\rho_{\text{dish}}$) of 1040 kg/m$^3$. The ultrasonic wave velocity ($c_{\text{dish}}$) was measured at 2367 m/s using the method described in [18]. The attenuation ($\alpha_{\text{dish}}$) of the ultrasonic waves in the material was about 1.1 dB/cm at 1 MHz, and 1.25 dB/cm at 2.25 MHz.

The anti-reflection cover was a custom-designed 38 mm-high thermoplastic polyoxymethylene (POM) tube containing degassed water. The wall thickness of the POM tube is 1 mm. A thin film of 40 µm, stretched at the bottom of the POM tube, ensured its insulation from the degassed water. The thin film material was compatible with any cell culture medium. To meet the objective of avoiding wave reflection and considering the constraints in dimensions of cell culture in an incubator (preventing addition of a high amount of water), a 10 mm-thick absorbing foam (Aptflex F28, Precision Acoustics Ltd, Dorchester, UK) was inserted inside the POM tube, 28 mm above the film. The anti-reflection cover was placed over the Petri dish.

In order to adapt the Petri dish/anti-reflection cover set to the in vitro ultrasound cell stimulation, the diameters of the Petri dish and of the anti-reflection cover were defined to make gas exchanges with the outside possible. Therefore, the diameters of the Petri dish and of the anti-reflection cover were 60 mm, larger than the diameter of the culture area (of hypothetical cells) of 35 mm.

Finally, the Petri dish/anti-reflection cover set was held above the transducer by means of a holder that was large enough not to interfere with the acoustic field generated by the transducers. The assembly ensures the horizontality of the active face of the transducer, the Petri dish and the anti-reflection cover.

4 Experimental set-up and configurations

The general synoptic diagram of the experiments is presented in Figure 2. The experiment room was air-conditioned and kept at a constant temperature. Experimental conditions did not change during the tests. The experiments in this study were performed successively using two immersion transducers (Imasonic, Voray-sur-l’Oignon, France) of center frequency 1 MHz and 2.25 MHz,
respectively. The active front diameter and active area were 13 mm and 1.33 cm² for the 1 MHz transducer, and were 9.5 mm and 0.71 cm² for the 2.25 MHz transducer. The acoustic impedance of the piezo-composite material (Z_{trans}) was 21 MRayls.

Each transducer was powered first by a pulse-receiver generator (Sofranel 5077 PR, Olympus, Waltham, MA 02453, USA) for the calibration of the experiments, and then, by a waveform generator (TGA 1241, Thurlby Thandar Instruments Limited, Huntingdon, UK) to study the performance of the anti-reflection cover under LIUS conditions. The radio-frequency signals (RF-signals) were conveyed from the 12-bit oscilloscope (Lecroy HDO 6104, Teledyne Inc., Thousand Oaks, CA USA) to a personal computer using a USB interface file transfer, and stored. The processing algorithms were implemented using Matlab (The MathWorks, Inc., Natick, MA, USA).

Acoustic pressure measurements were performed in transmission mode using a 0.5 mm needle hydrophone, a submersible wide-band amplifier and a converter-coupler (Precision Acoustics Ltd, Dorchester, UK). According to the technical data sheet provided by the manufacturer, the sensitivity of the needle hydrophone was 426 mV/MPa for the 1 MHz-transducer, and 422 mV/MPa for the 2.25 MHz-transducer.

The transducer and the hydrophone placed above were immersed in a large circular tank (3 m diameter and 1.25 m height) filled with degassed water (ρ_{water} = 1000 kg/m³). The temperature of the water in the tank (17.8 °C) was determined using a digital thermometer.

The ultrasonic incident wave beam was perpendicular to the water–air interface as well as to the water–dish interface. The ultrasonic wave was therefore transmitted through the Petri dish under normal incident conditions. Only pure compression waves were taken into account, and the shear waves propagating in the walls of the Petri dish were considered negligible. The ultrasonic wave velocities were taken to be constant and independent of the frequency (non-dispersive medium and material).

A micrometric electro-mechanical positioning scanner (precision of 0.1 mm) was used to locate the hydrophone in the 3D-space (x, y, z) from the front face to the far field of the transducer. The distance between the transducer and the water–air interface was defined by D. Its value was estimated by measuring the arrival time of the echo returned by the water–air interface. In all configurations, the hydrophone remained fully immersed. Figure 3 show the different experimental configurations selected for the acoustic measurements:

- **High water level and vertical position of the hydrophone, without (Fig. 3a) and with the Petri dish (no cover) (Fig. 3b):** This configuration means that there was a high column of water above the transducer (D = 700 mm). The wave reflected from the water–air interface arrived with a significant delay compared with the main incident wave, and was outside the signal recording range.
- **Low water level and horizontal position of the hydrophone, with the Petri dish, without (Fig. 3c) and with (Fig. 3d) the anti-reflection cover:** This configuration modeled a liquid level similar to the one used by biologists in cell culture dishes. The distance D was 46.3 mm.

### 5 Calibration of the acoustic set-up

For the calibration of the set-up, including the characterization of the acoustic fields of the transducers, and the hydrophone and the Petri dish locations, the chosen configuration was the one presented in Figure 4a. There was no cover over the Petri dish, and the hydrophone was in vertical position. The ultrasonic wave transmission mode was the short-pulse mode described as follows.

#### 5.1 The short-pulse mode

The transmitted input signal was an electric short-pulse signal delivered by the pulse-receiver generator. The pulse repetition frequency \( f_{PRF} = \frac{1}{t_{PRP}} \) was 100 Hz, and the gain was fixed at 10 dB. The frequency calibrations of the pulse-receiver generator were 1 MHz and 2.25 MHz, respectively. The mean ultrasonic wave velocity in the tank (\( c_{water} = 1474 \pm 1 \) m/s) was calculated using the arrival time variation measurements for two distances between each transducer and the water–air interface (Fig. 4a). The attenuation of the ultrasonic waves in water (\( \alpha_{water} \)) was 0.0022 dB/cm, whatever the frequency.
5.2 The acoustic fields and the focal zones

With the micrometric electro-mechanical positioning scanner, the peak acoustic pressure distribution was measured by scanning the $xz$-plane ($y = 0$), with $x$ ranging from $-7$ to $7$ mm, and $z$ ranging from $4$ to $45$ mm from the front face of the transducer. Figures 4b and 4c (Figs. 4d and 4e) show the axial distribution of the ultrasonic pressure delivered by the 1 MHz-transducer (2.25 MHz-transducer) along the beam axis ($z$). The depths marking the beginning of the focal zone (in the sense of the definitions of Bushberg et al. [19]) were $20 \pm 2$ mm for the 1 MHz-transducer and $25 \pm 2$ mm for the 2.25 MHz-transducer. The lateral apertures were 12 mm and 7 mm, respectively. The theoretical focal length for the 1 MHz-transducer (respectively, for the 2.25 MHz-transducer) was of $29$ mm (respectively, $31$ mm) at $-6$ dB, and the corresponding lateral resolution was $13$ mm (respectively, $9$ mm). The theoretical and experimental values are therefore in agreement.

5.3 Location of the Petri dish

The distance between the transducer and the Petri dish back wall was defined by $h$. The distance $h$ was chosen so that the Petri dish back wall was located on the maximum acoustic peak pressure of the acoustical fields, that is, $25$ mm when using the 1 MHz-transducer, and $36$ mm when using the 2.25 MHz-transducer. Its value was estimated using the pulse-echo mode by measuring the arrival time of the echo returned by the Petri dish back wall.

5.4 Locations of the hydrophone

The distance $z_f$ was defined as the distance between the transducer and the tip of the hydrophone, without or with the Petri dish and the anti-reflection cover (Fig. 3). The distance $z_f$ was $28$ mm with the 1 MHz-transducer, and was $39$ mm with the 2.25 MHz-transducer, and was constant for all measurement configurations. As shown in
Figures 3c and 3d, the hydrophone was moved from a vertical position (0°) to a horizontal position (90°). This last position allowed measurements in the Petri dish water layer, with and without the anti-reflection cover.

A 4 mm-deep notch was made in the lateral wall of the Petri dish. In vertical and horizontal positions, the tip of the hydrophone was located in the center of the Petri dish at the height $a = 2.2$ mm above the bottom. According to the technical data sheet provided by the manufacturer [20], the frequency-dependent directional response of the hydrophone is angularly wide enough to justify the comparison of the measurements when it is in the vertical position or the horizontal position. The relative amplitude between the two positions was $-5$ dB for the 1 MHz-transducer, and $-10$ dB for the 2.25 MHz-transducer. The RF-signals recorded for each position were therefore normalized with respect to these relative amplitudes. Table 1 summarizes all distances and heights for the experiments performed at 1 MHz and 2.25 MHz.

6 Validation of the anti-reflection cover under LIUS conditions

To test the effectiveness of the anti-reflection cover to meet the LIUS requirements, in particular to ensure that the delivered acoustic intensity level corresponds to a Spatial-Average Temporal-Average Intensity, $I_{SATA} (9.2)$ of 30 mW/cm², the ultrasonic wave transmission mode was the tone-burst mode described below.

6.1 The tone-burst mode

The temporal source waveforms delivered by the waveform generator were a 1 MHz-burst and a 2.25 MHz-burst signal, with a time duration of 200 μsec. The amplitude of the waveform generator was 9.5 V (peak-to-peak) for the 1 MHz-transducer and 10 V (peak-to-peak) for the 2.25 MHz-transducer (50 Ω). The pulse repetition period ($t_{PRP}$) was fixed at 1 ms, through a duty cycle (DC) of
20%. Figure 5 shows an example of 1 MHz-transmitted RF-signal over two periods, and recorded by the hydrophone at $z_f = 28$ mm, without Petri dish.

### 6.2 Characterization of the ultrasound intensity distributions under LIUS conditions

The temporal-peak intensity $I_{TP}$, temporal-average intensity $I_{TA}$, spatial-peak temporal-average intensity $I_{SPTA}$ and spatial-average temporal-average intensity $I_{SATA}$ [mW/cm$^2$] (9.2) were measured – the hydrophone being in the vertical position – on a total sounded area around 2.25 cm$^2$ ($[x, y] = [15$ mm, $15$ mm]) with a linear pitch of 0.25 mm, at $z_f = 28$ mm for the 1 MHz-transducer and $z_f = 39$ mm for the 2.25 MHz-transducer. Two sets of experiments were conducted for each transducer, without (Fig. 3a) and with (Fig. 3b) the Petri dish, with a high water level ($D = 700$ mm).

The intensity $I_{TA}$ (see Eq. (A.4)) was calculated from the measurement of the peak acoustic pressure $p_0$ and the intensity $I_{TP}$ (see Eq. (A.3)) following the algorithm proposed by Harris [21] and Preston [22]. $I_{SPTA}$ was the maximum value of the intensity spatial distribution $I_{TA}$ (see Eq. (A.5)), and $I_{SATA}$ was the spatial average of $I_{TA}$ (see Eq. (A.6)) calculated on the $-6$dB beam area.

Figure 5 shows the $I_{TA}$ distributions for the 1 MHz-transducer. Figure 7 shows the results for the 2.25 MHz-transducer. The circle in the image represents the active front face of the transducer (13 mm). The corresponding $I_{SPTA}$ and $I_{SATA}$ are summarized in Table 2.

This analysis makes it possible to verify that the experimental conditions of insonification in the high water level configuration correspond to the values commonly used in LIUS therapy. The $I_{SATA}$ is equal to about 30 mW/cm$^2$ (31.6 mW/cm$^2$ at 1 MHz, and 21.9 mW/cm$^2$ at 2.25 MHz). Note that the $I_{SATA}$ value for the commercial device is generally given with a tolerance of about 30%. The results obtained for the field distributions and for the measurements of $I_{SATA}$ are within this tolerance, and are consequently in accordance with the LIUS conditions.

### 6.3 The anti-reflection cover similar to a large water column

Figure 8 shows the comparison of transmitted signals between the transducer and the hydrophone in horizontal position for a low water level ($D = 46.3$ mm) without (Fig. 3c) and with (Fig. 3d) the anti-reflection cover, and in vertical position for a high water level ($D = 700$ mm) without the anti-reflection cover (Fig. 3b). Figure 9 shows

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Table 1. Table of distances and heights.

| Variable | Distance between | Value (mm) |
|----------|------------------|------------|
| $z_f$    | The transducer$^1$ and the hydrophone | 28 (1 MHz), 39 (2.25 MHz) |
| $D$ (high water level) | The transducer$^1$ and the water-air interface | 700 |
| $D$ (low water level) | The transducer$^1$ and the water-air interface | 46.3 |
| $h$      | The transducer$^1$ and the Petri dish$^2$ | 25 (1 MHz), 36 (2.25 MHz) |
| $a$      | The Petri dish$^3$ and the hydrophone | 2.2 |

$^1$ The active front face of the transducer.

$^2$ The back wall.

$^3$ The bottom of the Petri dish.
the same results for the 2.25 MHz-transducer. Table 3 summarizes the mean peak acoustic pressures measured for all configurations.

In the low water level configuration, the mean peak acoustic pressure was higher without the anti-reflection cover (149 ± 7 kPa at 1 MHz, and 183 ± 5 kPa at 2.25 MHz) than with the anti-reflection cover (105 ± 1 kPa at 1 MHz, and 83.5 ± 0.6 kPa at 2.25 MHz, cf. Tab. 3). Without the anti-reflection cover, the ultrasonic waves were strongly reflected at the water–air interface, which resulted in a dramatic increase in the acoustic pressure in the Petri dish water layer. With the anti-reflection cover, the ultrasonic waves were transmitted inside and partially absorbed by the absorbing foam that prevents multiple reflections by suppressing the water–air interface. The reduction in the mean peak acoustic pressure was 30% at 1 MHz and 54% at 2.25 MHz. The acoustic pressure measured using the hydrophone in the horizontal position, under the anti-reflection cover in the low water level configuration, was of the same order of magnitude as that measured when the hydrophone was in a vertical position, without the anti-reflection cover and in the high water level configuration (105 ± 3 kPa at 1 MHz, and 86.7 ± 2 kPa at 2.25 MHz, cf. Tab. 3). In summary, there were no more standing waves in the Petri dish water layer

Figure 6. Temporal-average intensity $I_{TA}$ [mW/cm$^2$] distribution without (left – Fig. 3a) and with (right – Fig. 3b) the Petri dish for 1 MHz-transducer at $z_f = 28$ mm. The white circle represents the area of the transducer.

Figure 7. Ultrasound temporal-average intensity distribution $I_{TA}$ [mW/cm$^2$] without (left – Fig. 3a) and with (right – Fig. 3b) the Petri dish for 2.25 MHz-transducer at $z_f = 39$ mm. The white circle represents the area of the transducer.

Table 2. Measured peak acoustic pressure and intensities.

| Frequency/Configuration | $p_0$ [kPa] | $I_{SPTA}$ [mW/cm$^2$] | $I_{START}$ [mW/cm$^2$] |
|-------------------------|-------------|-------------------------|-------------------------|
| 1 MHz                   |             |                         |                         |
| Without Petri dish      | 103         | 74.4                    | 37.7                    |
| With Petri dish         | 100.5       | 60.7                    | 31.6                    |
| 2.25 MHz                |             |                         |                         |
| Without Petri dish      | 118.5       | 80.7                    | 42.4                    |
| With Petri dish         | 93.5        | 51.4                    | 21.9                    |

*–6 dB.
when the anti-reflection cover was present, and the intensity delivered was correctly estimated by measuring the mean peak acoustic pressure in the Petri dish water layer with a high water column above. It is worth noting that the presence of the hydrophone in the horizontal position necessarily disturbs the acoustic field. However, as it is present in all configurations, this disturbance is identical, even though, for the configuration with the anti-reflection cover and the 40 µm-film, it could disturb more. The presence of the hydrophone could also explain the high amplitude of

**Figure 8.** Moduli of the Hilbert transform of the transmitted RF-signals between the 1 MHz-transducer and the hydrophone in horizontal position and a low water level above the Petri dish, and the hydrophone in vertical position and a high water level. On the right, the mean values and standard deviations are mentioned.

**Figure 9.** Moduli of the Hilbert transform of the transmitted RF-signals between the 2.25 MHz-transducer and the hydrophone in horizontal position and a low water level above the Petri dish, and the hydrophone in vertical position and a high water level. On the right, the mean values and standard deviations are mentioned.

**Table 3.** Mean peak acoustic pressure measured, $p_0$ [kPa].

| Frequency MHz | Low water level without cover | Low water level with cover | High water level without cover |
|---------------|-------------------------------|----------------------------|--------------------------------|
| 1             | 149 ± 7                       | 105 ± 1                    | 105 ± 3                        |
| 2.25          | 183 ± 5                       | 83.5 ± 0.6                 | 86.7 ± 2                       |
the signals, standing waves being created under the hydrophone. To conclude, the effect of the anti-reflection cover is similar to that of a high water column above the Petri dish, reducing drastically the perturbation induced by the wave reflection at the water–air interface.

### 6.4 Analytical estimation of the acoustic intensity under the anti-reflection cover

Once the effectiveness of the anti-reflection cover was validated, a quantitative study of the energy transmitted into the water layer was undertaken, and the intensity loss calculated (see Eq. (A.8)). To do so, a comparison between the analytically calculated $I_{\text{SPTA}}^\text{A}$ and measured $I_{\text{SPTA}}^\text{M}$ intensities was conducted. The time duration of the transmitted signal was 200 μsec. Therefore, the propagation distance was about 29.5 cm, which was of the same order of magnitude as the distance between the transducer and the Petri dish back wall ($h = 25$ mm, and $h = 36$ mm, Tab. 1). $I_{\text{SPTA}}^\text{M}$ measured at $a = 2.2$ mm from the bottom of the Petri dish was the result of a sum of several waves that traveled back and forth several times between the transducer and the Petri dish back wall, and inside the bottom thickness of the Petri dish. Therefore, for the analytical calculation, these back-and-forth travels (one Run is a one-way path traveled by the ultrasonic wave) were considered. Three runs were reasonable, given the distances between the interfaces. Multiple reflections beyond three runs were not significant. Recordings in reflection and pulse-echo mode (back wall echoes on the Petri dish) were made to analyze the amplitudes of the RF-signals due to the multiple reflections, and beyond four back-and-forth travels ($3 + 1$ Runs) the signals were weak, and slightly out of the noise.

The reflection and transmission coefficients (see Eqs. (A.1) and (A.2)) were calculated at the water–Petri dish interface, at the water–transducer interface, and inside the bottom thickness. The attenuations (see Eq. (A.7)) were taken into account in the water and in the thickness of the Petri dish. The parameters of the study are summarized in Table 4.

The loss of intensity of a wave (see Eq. (A.8)) inside the Petri dish, at a distance $z_I$ along the beam axis of the transducer is presented in Table 5 for the 1 MHz-transducer, and in Table 6 for the 2.25 MHz-transducer. The intensity increases under the Petri dish because of the multiple reflections between the dish and the transducer.

Analytical and measured intensities are similar for both transducers. The differences between the two values are 2.5% (1 MHz) and 6% (2.25 MHz). This result is acceptable and a simple analytical calculation makes it possible to estimate the ultrasound intensity the cells in the culture medium receive.

### 7 Discussion & conclusion

Low intensity ultrasound stimulation is a modality used today in medical therapy for bone regeneration. The performance of the commercial devices used in the clinical field is still being debated. To understand the physical phenomena, laboratory tests are performed on cells placed in Petri dishes and insonified by acoustic waves. Unfortunately, as mentioned in [13], several parameters are liable to compromise the monitoring of the acoustic dose delivered to the cells. Among them, the distance between the transducer and the Petri dish, the coupling medium (degassed water), and the formation of standing waves inside the culture medium. All of them have been treated in the anti-reflection cover device. A fourth one is the relative size of the beam and the dish and the homogeneity of the field delivered inside the Petri dish. This question is in progress. In this work, a custom-designed anti-reflection cover placed on the Petri dish has been developed, whose constituent materials and dimensions are entirely compatible with in vitro cell tests in incubators. This anti-reflection cover makes it possible to trap transmitted waves and to avoid the standing waves in the cell culture medium. Therefore, the cell culture medium is comparable to a column of water several centimeters high in which the waves would not be reflected on any interface. Even if the results presented in this paper (Figs. 8 and 9) show the efficiency of the cover, more tests are needed to improve the reproducibility of these results and to obtain a fine interpretation of the interaction of the ultrasonic field with the anti-reflection cover. On this last point, the development of a numerical model reproducing the experimental conditions is in progress and should bring some answers. Thus, the acoustic intensity in the cell culture medium inside the Petri dish/anti-reflection cover set can be estimated by the measurement of an equivalent intensity in a large volume of water above a Petri dish. It is then possible to estimate the energy received by cells under low intensity ultrasound stimulation. This set-up has been designed (dimensions and materials) to allow in vitro ultrasound stimulation without having to remove the cell cultures from the incubator. It offers the possibility

| Physical parameters | Values |
|---------------------|--------|
| $c_{\text{water}}$ | $1474 \pm 1$ m/s |
| $\rho_{\text{water}}$ | 1000 kg/m$^3$ |
| $Z_{\text{water}}$ | 1.47 MRays |
| $\tau_{\text{water}}$ | 0.0022 dB/cm |
| $c_{\text{dish}}$ | 2367 m/s |
| $\rho_{\text{dish}}$ | 1040 kg/m$^3$ |
| $Z_{\text{dish}}$ | 2.46 MRays |
| $\tau_{\text{dish}}$ | 1.10 dB/cm (1 MHz) |
| $E^*$ | 1.25 dB/cm (2.25 MHz) |
| $a^{**}$ | 0.8 mm |
| $R_{\text{water,dish}}$ | 2.2 mm |
| $Z_{\text{trans}}$ | 0.063 |
| $Z_{\text{water,transducer}}$ | 21 MRays |
| $R_{\text{water,transducer}}$ | 0.755 |

* The bottom thickness of the Petri dish.
** The distance between the bottom of the Petri dish and the hydrophone tip position.
to avoid any human manipulation of the cell cultures even in the case of an ultrasound stimulation practiced over several days which is excluded in the configurations proposed in [16] and [10]. This possibility greatly limits the risks of contamination of cell cultures. Moreover the anti-reflection cover is easily adaptable to any current well on water surface whatever its dimensions and even for multwell plates. Its low cost makes it possible to consider duplicating it for serial studies and achieve the amount of data necessary for statistical analysis of biological results.

Conflict of interest

The authors declare that they have no conflicts of interest in relation to this article.

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Table 5. Analytical calculation of the intensity and loss for the 1 MHz-transducer.

| Intensity of the wave                                      | (mW/cm²) | Intensity loss (%) |
|-----------------------------------------------------------|----------|--------------------|
| Measured at 25 mm (under the Petri dish)                  | 71.6     | 15.2               |
| Measured at 28 mm (above the Petri dish)                  | 60.7     | 11.9               |
| Calculated after 3 runs in water                          | 78.4     | 3.2                |
| Transmitted into the Petri dish                           | 73.5     | 5.2                |
| Calculated after 3 runs in the bottom thickness           | 62.3     | 13                 |

Table 6. Analytical calculation of the intensity and loss for the 2.25 MHz-transducer.

| Intensity of the wave                                      | (mW/cm²) | Intensity loss (%) |
|-----------------------------------------------------------|----------|--------------------|
| Measured at 36 mm (under the Petri dish)                  | 64.5     | 20.3               |
| Measured at 39 mm (above the Petri dish)                  | 51.4     | 20.3               |
| Calculated after 3 runs in water                          | 70.6     | 11.9               |
| Transmitted into the Petri dish                           | 66.1     | 13                 |
| Calculated after 3 runs in the bottom thickness           | 54.7     | 15.2               |
A Wave coefficients and intensities

A.1 Reflection and transmission coefficients

The reflection coefficient of a wave on a medium 1/medium 2 interface is defined under normal incident conditions as:

\[ R_{1,2} = \frac{Z_2 - Z_1}{Z_2 + Z_1}, \]

and the transmission coefficient is defined as:

\[ T_{1,2} = \frac{4Z_1Z_2}{Z_2 + Z_1}, \]

where \( Z_1 = \rho_1c_1 \) and \( Z_2 = \rho_2c_2 \) are the acoustic impedances [MRayls] of medium 1 and of medium 2. \( \rho \) is the mass density [kg/m\(^3\)] of the medium, and \( c \) is the ultrasonic wave velocity [m/s] propagating in the medium.

A.2 Acoustic intensities

The methods for assessing acoustic intensities were proposed by Harris [21] and Presten [22]. There are several definitions of useful acoustic intensities referenced by international authorities [23], including the temporal-peak intensity \( I_{TP} \) [mW/cm\(^2\)] which describes the distribution of the acoustic pressure of a wave over time, at the position \( z \) along the beam axis, and which is calculated as:

\[ I_{TP}(z) = \frac{p_0(z)^2}{Z} \leq \frac{\max[p^+(z),p^-(z)]^2}{Z}, \]

where \( Z \) is the acoustic impedance of the medium. The peak acoustic pressure \( p_0 \) [kPa] is therefore assumed to be the greater of the two values between the maximum peak acoustic pressure \( p^+ \) (i.e., maximum voltage of the recorded signal) and the minimum peak acoustic pressure \( p^- \) (i.e., minimum voltage of the recorded signal). The Temporal-Average Intensity \( I_{TA} \) [mW/cm\(^2\)] is the intensity \( I_{TP} \) over the total time of the ultrasonic insonification, equal to the ratio between the pulse/burst duration \( t_{pulse} \) and the pulse repetition period \( t_{PRP} \):

\[ I_{TA}(z) = I_{TP}(z) \frac{t_{pulse}}{t_{PRP}}. \]
By cross-referencing spatial and temporal considerations, several intensities can be defined. The spatial-peak temporal-average intensity \( I_{SP\text{T}A} \) [mW/cm\(^2\)] is the maximum intensity in the spatial field cross-section at a distance \( z_f \) from the front face of the transducer:

\[
I_{SP\text{T}A}(z_f) = \max_{x,y}(I_{TA}(z_f)),
\]

where \( x \) and \( y \) are the coordinates of the points on the surface.

The spatial-average temporal-average intensity \( I_{SATA} \) [mW/cm\(^2\)] is the measure of intensity \( I_{TA} \) per unit area at the distance \( z_f \) from the front face of the transducer:

\[
I_{SATA}(z_f) = \frac{1}{N} \sum_{x,y=1}^{N} [I_{TA}(z_f)]_{x,y},
\]

\[
\geq 0.25 \times I_{SP\text{T}A}(z_f),
\]

(A.5)

(A.6)

A.3 Intensity loss

The intensity \( I(z_2) \) of a wave at the distance \( z_2 \) is function of the initial intensity \( I(z_1) \) and the attenuation coefficient \( a \) of the wave over the distance \( |z_2 - z_1| \):

\[
I(z_2) = I(z_1) \exp(-2a|z_2 - z_1|).
\]

(A.7)

The intensity loss is expressed as a percentage (%) of the intensity lost:

\[
I_{\text{loss}} = 100 \times \frac{I(z_1) - I(z_2)}{I(z_1)}.
\]

(A.8)

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