En-face Tomographic Imaging of Scattering Objects Using Single Broadband Light Emitting Diode Based Full-Field Optical Coherence Microscopy

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Structural and morphological elucidation of soft scattering objects such as biological tissues is an important aspect of study for the development of diagnostic tools. Optical methods suffer from the limitation of scattering and absorption of light in such kind of objects; limiting the applicability in deep tissue characterizations. In this present work, we used a high-resolution full-field optical coherence microscopy (FF-OCM) with a single broadband light emitting diode (470–850 nm) as illumination source for the imaging of scattering biological objects. The FF-OCM system was based on the Linnik geometry and a two-dimensional charge complementary oxide semiconductor camera (CMOS). The sequential two-dimensional spatial multiple phase-shifted interferograms were obtained by moving the sample stage using a piezoelectric transducer and the en-face FF-OCM images were subsequently reconstructed by a fast derivative-based method. The spatial resolution of the present system was 0.9 (air) and 1.38 µm axially and laterally, respectively.

The present system was demonstrated for the imaging highly scattering artificial skin dressing, onion, and fish skin samples. Our results show the capability of the present system to clearly discern the structural features of the fish skin as evident from the depth intensity profiles at different positions. The present system is significantly stable, compact, and cost-effective with comparable spatial resolution compared to conventional FF-OCM systems and can be further used to distinguish a normal tissue from a diseased tissue.

Key words: full-field optical coherence microscopy; light emitting diodes; derivative-based method; scattering objects.

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Introduction

Optical imaging techniques have revolutionized the biomedical research. These imaging techniques provide several advantages such as non-invasive, cost-efficient, fast, and pose no health side effects compared to imaging modalities using ionizing radiations. Traditional optical methods can provide micron-scale resolution at the sub-cellular and sub-organelle level at superficial depths of tissues. However, to see deep inside intact biological tissues, these light-based imaging modalities suffer from several challenges due to the high scattering and absorption of light [1, 2]. But, statistical characteristics of scattering or attenuation coefficient of light can provide structural/morphological information at micrometer, or submicrometer resolution for disease diagnosis at the early phase [3]. Thus, imaging modalities which can explore this scattering of light as contrast for image reconstruction can provide information from deep depth of such specimens. Optical coherence tomography (OCT), as demonstrated by James Fujimoto at 1991, is such a technique and based low coherence interferometry [4]. Moreover, optical coherence microscopy (OCM) uses higher numerical aperture objectives for non-invasive, micron-scale cross-sectional depth-resolved imaging. OCT/OCM, till date, has been applied for several biomedical and non-destructive testing of materials [5].
Among the several types of OCT, full-field OCT/OCM (FF-OCT/FF-OCM) can provide high-resolution images of intact tissues to characterize their structural features [6]. FF-OCM can provide XY en-face sectional images in a single exposure using two-dimensional (2D) charge coupled device [6], or charge complementary oxide semiconductor (CMOS) cameras [7]. The en-face sectional imaging is complementary to the longitudinal sectional imaging as it can provide temporal information within the en-face plane. Various FF-OCT/FF-OCM systems based on both time and frequency domain techniques for the quantitative imaging of biological samples have been developed [6–10]. In general, broadband infrared superluminescent diodes (SLDs) and white light sources based on thermal, halogen and xenon arc lamp have been used as illumination sources in FF-OCM to achieve high axial resolution [10–12]. High power halogen light sources have advantages as it provides highest axial resolution in FF-OCM till date [13].

Among the several possible light sources, as mentioned above, light emitting diodes (LEDs), compact in design and cost-effective; have also been used as alternative interferometric illumination sources [14, 15]. LEDs, compared to conventional coherent light sources, provide significant advantages due to its short temporal coherence and spatial incoherence which eliminate the background and speckle noises. LEDs have been used, due to its advantages, in white light interferometry, digital holography, and vertical scanning interferometry [16–18]. A FF-OCT system with an infrared LED has been reported, providing higher resolution compared to SLD-based systems [4]. Infrared LED-based (central wavelength ($\lambda_c$) ~850 nm) FF-OCM systems have also been used for non-destructive testing of polymer coating on pharmaceutical pellets [19], skin imaging (with multi-wavelength RGB LED) [20] etc. A FF-OCM system with nitride-based white light emitting diode (WLED) ($\lambda_c$ ~532 nm; full-width-half-maxima ~190 nm) has also been reported for spectroscopic studies [21]. Till recently, Anna et al. [22] demonstrated the capability of WLED-based FF-OCM system for the characterizations of melanoma mice ear skin. WLED bears advantages over other broadband light sources such as halogen light due to its higher energy efficiency, longer lifetime, compact, less heat dissipation, and can provide high axial resolution.

In this work, we demonstrated the capability of a single broadband LED-based (470–850 nm, $\lambda_c$=630 nm) Linnik type FF-OCM system to study highly scattering samples. Scattering can provide an endogenous contrast for the center wavelength of the light source i.e., $\lambda_c$, $\psi(x, y)$ is the phase delay due to the sample and $\gamma(x, y, t)$ is the envelope of autocorrelation of the light source. When the phase shift between two sequential images is very small compared to coherence length of the light source, then $\gamma(x, y, t)=1$. However, multiple phase-shifted images were taken in a time sequence. Subsequently first, second, and third derivatives of $l(x, y, t)$ with respect to scanning time $(t)$ was expressed in terms of discrete differences of four sequential images i.e., $l(x, y)$, $l(x, y)$, $l(x, y)$, and $l(x, y)$. Combining these images, the en-face sectional amplitude images at $p^\text{th}$ layer without the requirement of the phase-shift angle can be obtained as [22, 23]:

$$l(x, y, t)=l_0(x, y)+\gamma(x, y, t)A(x, y, t)\sin(\omega t)\psi(x, y),$$

(1)

where $l_0(x, y)$ is the background intensity, $\omega=2\pi f$ is the central frequency of the light source corresponds to the center wavelength of the light source i.e., $\lambda_c$, $\psi(x, y)$ is the phase delay due to the sample and $\gamma(x, y, t)$ is the envelope of autocorrelation of the light source. When the phase shift between two sequential images is very small compared to coherence length of the light source, then $\gamma(x, y, t)=1$. However, multiple phase-shifted images were taken in a time sequence. Subsequently first, second, and third derivatives of $l(x, y, t)$ with respect to scanning time $(t)$ was expressed in terms of discrete differences of four sequential images i.e., $l(x, y)$, $l(x, y)$, $l(x, y)$, and $l(x, y)$. Combining these images, the en-face sectional amplitude images at $p^\text{th}$ layer without the requirement of the phase-shift angle can be obtained as [22, 23]:

$$A(x, y)=\left[\{l_1(x, y)-2l_0(x, y)+l_2(x, y)\}^2-\{l_0(x, y)-3\{l_2(x, y)-l_1(x, y)\}+l_0(x, y)\}l_0(x, y)\right]^{1/2}.$$
Figure 1. Schematic of the single broadband light emitting diode based full-field optical coherence microscopy system
LED: light emitting diode; L₁–L₃: lenses; CMOS: complementary oxide semiconductor; PC: personal computer; PZT: piezoelectric transducer

Figure 2. Characterization of the single broadband light emitting diode based full-field optical coherence microscopy system:
(a) response of the coherence function; (b) en-face OCT image of USAF 1951 test target; (c) intensity profile of the eighth group taken at the place indicated by the yellow line in (b)

The performance of the single broadband LED-based FF-OCM system for the imaging of scattering objects was evaluated by measuring the spatial (axial and lateral) resolution and signal to noise ratio (SNR). The axial resolution of the FF-OCM system is governed by the coherence length of the light source. Figure 2 (a) shows the measured coherence function as 0.9 µm by using a plane reflecting mirror as a sample [22]. The axial resolution using full-width-half maxima, Δλ=212 nm (measured) and refractive index n=1 (air) was calculated as 0.82 µm; thus showing both the calculated and measured axial resolutions were in good agreement with each other. In
addition, the SNR was also measured from the obtained axial response. The measurement was performed by inserting an optical density filter (~1.3) in the sample arm with a round trip attenuation of 26 dB. A variable attenuator was used in the reference arm to keep CMOS close to the saturation level. The SNR of the system was obtained to be ~44 dB and limited by the full well pixel capacity of the CMOS camera. Further, the lateral resolution was also measured by reconstructing the en-face image of a standard USAF 1951 test target. The intensity profile across the lines of group 8 was plotted, and the lateral resolution of the system was obtained to be ~1.38 µm. The calculated (~1.32 µm) and the measured lateral resolutions showed similar performance of the system.

We also characterized our single broadband LED-based FF-OCM system in detail by measuring the light source stability in terms of noise to signal ratio i.e., 0.50% without averaging and 0.20% with 3×3 pixel averaging at nearby saturation and phase stability in terms of standard deviation in phase i.e. 11 mrad. The detail measurement processes of these parameters also can be found in our previous research article by Anna et al. [22, 24].

Results and Discussion

Here, we imaged several scattering objects using our developed single broadband LED-based FF-OCM system such as, an artificial skin dressing, onion, and fish skins. The results of our measurements and the corresponding discussion are mentioned as follows.

Figure 3. Single broadband-based full-field optical coherence microscopy imaging of an artificial skin dressing: (a) photograph of skin dressing sample (indicating imaging areas; position 1 and 2); (b)–(e) en-face sectional images of skin dressing sample at 15, 38, 60, and 90 µm corresponds to position 1 and (f)–(i) at 15, 35, 55, and 95 µm depths corresponds to position 2, respectively; (j) and (k) shows stacked 3D volume (500×500×100 µm) images corresponding to positions 1 and 2, respectively as marked in the photograph of the sample (a). Scale bar: 50 µm
At first, the developed FF-OCM system was used to study a highly scattering artificial skin dressing sample (Skin Temp; Taiyan Biotech Co., Ltd., China) was chosen. Skin dressings are materials which are generally used to protect wounds during healing process. These are multilayered composite structures, and can serve as an ideal scattering sample, mimicking very close to real biological tissues. The incident power of the light at the sample plane and the exposure time were 60 µW and 10 ms, respectively. The photograph of the highly scattering complex structured artificial skin dressing sample is shown in Figure 3 (a). The imaging areas are marked by position 1 and 2. The sequentially phase-shifted interferograms of the artificial skin dressing at different depths were acquired (50 fps) by moving the PZT with constant step of 33 nm. To achieve higher depth in the sample, five sequential scans, each of 25 µm, by moving the PZT was performed. This was achieved by a computer controlled translation stage attached to sample stage (see Figure 1). A series of *en-face* FF-OCM images were then reconstructed using derivative-based method (using Eq. (2)) from the top of the sample surface to the different depths. Figure 3 (b)–(e) shows the *en-face* FF-OCM images at 15, 38, 60, and 90 µm depths for position 1 and (f)–(i), at 15, 35, 55, and 95 µm depths for position 2 of the sample, respectively, revealing the different complex structures of the sample. The 3D data set was obtained by combining the *en-face* sectional images and the 3D projection of volume 500×500×100 µm for position 1 and 2 are shown in Figure 3 (j) and (k), respectively, discerning the layered structure of the sample.

To further evaluate the capability of imaging of scattering biological samples by the developed single broadband-based FF-OCM system, we measured the *en-face* images of an onion bulk and fish skin samples. The multiple phase-shifted interferograms were measured as mentioned previously and the *en-face* images at different depth were reconstructed. Figure 4 (a)–(c) shows the *en-face* FF-OCM images at 2, 20, 40, and 64 µm depths, respectively. Figure 4 (d) shows the stacked 3D volume image of 500×500×90 µm. The different cell organizations such as a cell wall inside the cell (epithelium) as well as concentric laminae layers of the onion skin are clearly resolved in these sectional images. Our approach can be implemented in studying physiological and/or pathological defects in onions and possibly other food products; thus can play an important role in food security.

We have also taken the Japanese butterfish (*Psenopsis anomala*) skin (Figure 5 (a)) as an example to observe the performance of the present system in elucidating structural complexity. The *en-face* FF-OCM images at 5, 18, 34, and 45 µm depths are shown in Figure 5 (b)–(e), which reveals structural heterogeneity at different depths of epidermis layer of the skin. The scales present in the upper layers of the skin (Figure 5 (b)) and complex structure in between the scales and joints (Figure 5 (c)) are clearly visible. At 34 and 45 µm depths, microridges of the epidermis can also be seen in Figure 5 (d) and (e). Figure 5 (f) shows the cross-sectional YZ view of the sample revealing the upper epidermis scale and stratum spongiosum with scale layers of epidermis structure of the fish skin. Figure 5 (g) shows the 3D volume image of 500×500×120 µm depicting this structural information, which can be used for further understanding of fish skin non-invasively. The residual fringes at higher depths may be attributed to the irregular pattern of the skin surface, arising probably due phase error occurrence.

Furthermore, the depth profile along the Z-direction for the fish skin sample was also plotted for better understanding of the sample structure. The depth profile of the FF-OCM intensity for 15×15 pixels for three different positions is shown in Figure 6 corresponding to the area marked by yellow boxes in the *en-face* sectional image as indicated by 1, 2, and 3 in Figure 5 (c).

Figure 6 (a) and (b) shows three prominent peaks in the depth profile and some of the *en-face* FF-OCM images can be seen in Figure 5 (c)–(e) corresponding to first two peaks. These peaks can be attributed to fish skin epidermis scale layer and stratum spongiosum.
Figure 5. En-face sectional imaging of Japanese butterfish skin using single broadband-based full-field optical coherence microscopy:
(a) photograph of the fish (indicating imaging area; red rectangular box); (b)–(e) en-face sectional images of fish skin at 5, 18, 34, and 45 µm depths, respectively; (f) cross-sectional view along YZ; (g) stacked 3D volume image of 500×500×120 µm. Scale bar: 50 µm

Figure 6. Depth intensity profiles obtained by single broadband-based full-field optical coherence microscopy:
(a)–(c) the depth profile of FF-OCM intensity corresponding to areas 1, 2, and 3 (15×15 pixels) marked in the en-face images of Figure 5 (c), respectively
region. Figure 6 (c) shows only one prominent peak in-depth profile which may be attributed to the first surface of epidermis with remaining scale part. Due to the high reflectivity from this region it is not possible to see deep inside. We also applied the polyfitting to these depth intensity profiles, showing distinctly different attenuation level of incident light. The slopes of the fitted lines for the intensity profiles in Figure 6 (a)-(c) are 0.0046, 0.0047, and 0.0056, respectively. As evident from the intensity profiles and fitted slope values, it can be concluded that our FF-OCM system provides layered structural information of the fish skin. Moreover, our results indicate the high stability of the present system to measure structural features of biological samples. The present system also offers higher spatial resolutions compared to SLD-based systems and comparable with other popular single broadband light source such as halogen. In addition, single broadband LEDs provide uniform illumination of the sample, which minimizes the background noise as compared to SLD-based FF-OCM.

Conclusion

In conclusion, a single broadband LED-based FF-OCM system was used to study structural features of highly scattering samples. As examples, artificial skin dressing, onion, and fish skins were studied. Our system provides high spatial resolutions viz. 0.9 µm (axial) and 1.38 µm (lateral) with measured sensitivity –44 dB. For interferometric image reconstruction, a derivative-based method was used. Our results show the capability of the present system in elucidating the structural features of our chosen scattering samples. Intensity profiles along depth in fish skin provided information of the skin layer organization. Moreover, different cellular structures can be seen in onion skin samples. These findings confirm the robustness of the present system to further its application in human skin (dermatology) study and microstructural assessment of normal and diseased tissues.

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Conflicts of interest. The authors have no conflict of interests to disclose.

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