SUBSTELLAR-MASS COMPANIONS TO THE K-DWARF BD+14 4559 AND THE K-GIANTS HD 240210 AND BD+20 2457

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ABSTRACT

We present the discovery of substellar-mass companions to three stars by the ongoing Penn State–Toruń Planet Search conducted with the 9.2 m Hobby–Eberly Telescope. The K2-dwarf, BD+14 4559, has a 1.5 \( M_J \) minimum mass companion with the orbital period of 269 days and shows a non-linear, long-term radial velocity (RV) trend, which indicates a possible presence of another planet-mass body in the system. The K3-giant, HD 240210, exhibits RV variations that require modeling with multiple orbits, but the available data are not yet sufficient to do it unambiguously. A tentative, one-planet model calls for a 5.2 \( M_J \) minimum mass planet in a 502 day orbit around the star. The most massive of the three stars, the K2-giant, BD+20 2457, whose estimated mass is 2.8 \( \pm 1.5 M_\odot \), has two companions with the respective minimum masses of 21.4 \( M_J \) and 12.5 \( M_J \) and orbital periods of 380 and 622 days. Depending on the unknown inclinations of the orbits, the currently very uncertain mass of the star, and the dynamical properties of the system, it may represent the first detection of two brown dwarf-mass companions orbiting a giant. The existence of such objects will have consequences for the interpretation of the so-called brown dwarf desert known to exist in the case of solar-mass stars.

Key words: planetary systems – stars: individual (BD+14 4559, HD 240210, BD+20 2457) – stars: low-mass, brown dwarfs

1. INTRODUCTION

Searches for planets around giant stars offer an efficient way to extend studies of planetary system formation and evolution to stellar masses substantially larger than 1 \( M_\odot \). (Sato et al. 2003, 2008a, 2008b; Hatzes et al. 2006; Hekker et al. 2008; Döllinger et al. 2009; Niedzielski et al. 2007, 2009; Liu et al. 2009). Although searches for the most massive substellar companions to early-type stars are possible (Galland et al. 2005; Lagrange et al. 2009), it is much more efficient to utilize the power of the radial velocity (RV) method by exploiting the fact the GK-giants, the descendants of the main sequence A–F type stars, have cool atmospheres and sufficiently many narrow spectral lines to make achieving a < 10 m s\(^{-1}\) RV measurement precision possible.

The GK-giant surveys are beginning to provide the statistics, which are needed to constrain the efficiency of planet formation as a function of stellar mass and chemical composition. In fact, the initial analyses by Johnson et al. (2007) and Lovis & Mayor (2007) extend to giants the correlation between planetary masses and the masses of their primaries observed for the lower mass stars. Most likely, this is simply because massive stars tend to have more massive disks. These results are in accord with the core accretion scenario of planet formation (Kennedy & Kenyon 2008). Furthermore, Pasquini et al. (2008) have used the apparent lack of correlation between the frequency of planets around giants and stellar metallicity to argue that this effect may imply a pollution origin of the observed planet frequency–metallicity correlation for main-sequence stars (Fischer & Valenti 2005). Similar conclusions have been reached in a preliminary study of our giant star sample by Zieliński et al. (2009).

The steadily extending baselines of the ongoing surveys of post-MS giants begin furnishing multiplanet system detections that are needed to study the dynamical evolution of planets around the off main-sequence stars. The multiplicity of planetary systems around MS stars has been firmly established (Udry & Santos 2007; Wright et al. 2009) and there is no reason why this tendency should not be present among giants. In fact, the ongoing search for planets around GK-giants with the Hobby–Eberly Telescope (HET) by our group (Penn State–Toruń Planet Search, hereafter PTPS) has already tentatively identified a two-planet system around the K0-giant HD 102272 (Niedzielski et al. 2009), and long-term trends in the RV data have been seen in other surveys (e.g., Sato et al. 2008a).

In this paper, we describe the detection of a planet around the K3-giant, HD 240210, and show evidence that the star has more low-mass companions, and the intriguing discovery of two brown dwarf-mass bodies orbiting the K2-giant, BD+20 2457. We also describe the discovery of a Jupiter-mass planet and a non-linear RV trend in the K2-dwarf, BD+14 4559. This detection is the first result of the extension of the PTPS to evolved dwarfs in the upper envelope of the MS. This part of the project will be described in detail in a future paper.

The plan of this paper is as follows. An outline of the observing procedure and a description of the basic properties of the three stars are given in Section 2, followed by the analysis of RV measurements in Section 3. The accompanying analysis of rotation and stellar activity indicators is given in Section 4. Finally, our results are summarized and further discussed in Section 5.

2. OBSERVATIONS AND PROPERTIES OF THE STARS

Observations were made with the HET (Ramsey et al. 1998) equipped with the High Resolution Spectrograph (HRS; Tull 1998) in the queue scheduled mode (Shectone et al. 2007). The spectrograph was used in the \( R = 60,000 \) resolution mode with a gas cell (\( I_2 \)) inserted into the optical path, and it was fed with a
Spectral type K2V K3III K2II analysis have been described in detail elsewhere (Niedzielski after Z09), who have estimated their values using the method of the K0 giant BD+70 1068 ($\eta$ $\pm$ $\delta$ $\epsilon$ $\zeta$).

Figure 1. Relative RVs for K0-giant BD+70 1068. The star shows RV variability at $\sigma$ = 12 m s$^{-1}$ whereas the intrinsic RV measurement uncertainty for these observations was 7 m s$^{-1}$. The remaining RV scatter can be attributed to solar-type oscillations at the level of 10 m s$^{-1}$.

Relative RVs of BD+14 4559, HD 240210 and BD+20 2457

| Parameter | BD+14 4559 | HD 240210 | BD+20 2457 |
|-----------|------------|-----------|------------|
| V         | 9.63       | 8.33      | 9.75       |
| $B-V$     | 0.98 ± 0.08| 1.65 ± 0.04| 1.25 ± 0.09|
| Spectral type | K2V | K3III | K2II |
| $\pi$ (mas) | 19.99 ± 1.61 | 7.0 ± 2.6 | 5.0 ± 26 |
| $T_{\text{eff}}$ (K) | 5008 ± 20 | 4297 ± 25 | 4137 ± 10 |
| log g     | 4.60 ± 0.05 | 2.31 ± 0.11 | 1.51 ± 0.05 |
| [Fe/H]    | 0.10 ± 0.07 | −0.18 ± 0.12 | −1.00 ± 0.07 |
| log $L/L_\odot$ | −0.32 | 1.55 | 3.17 |
| $M_\star/M_\odot$ | 0.86 ± 0.15 | 0.82 ± 0.15 | 2.8 ± 1.5 |

2" fiber. Details of our survey, the observing procedure, and data analysis have been described in detail elsewhere (Niedzielski et al. 2007; Niedzielski & Wolszczan 2008).

RVs were measured using the standard I$_2$ cell calibration technique (Butler et al. 1996). A template spectrum was constructed from a high-resolution Fourier transform spectrometer (FTS) I$_2$ spectrum and a high signal-to-noise stellar spectrum measured without the I$_2$ cell. Doppler shifts were derived from least-square fits of template spectra to stellar spectra with the imprinted I$_2$ absorption lines. The RV for each epoch was derived as a mean value of the independent measurements from the 17 usable echelle orders with a typical, intrinsic uncertainty of 6–8 m s$^{-1}$ at 1$\sigma$ level. This RV precision levels made it quite sufficient to use the Stumpff (1980) algorithm to refer the measured RVs to the Solar System barycenter.

The long-term precision of our RV measurements has been verified by the analysis of data derived from the monitoring of stars that do not exhibit detectable RV variations. The results for the KO giant BD+70 1068 ($V = 9.29$, $B-V = 1.15$ ± 0.04, $T_{\text{eff}} = 4580 \pm 17$ K) are shown in Figure 1 as an example. The relative RVs for this star vary with $\pi = 12$ m s$^{-1}$, whereas our estimated intrinsic RV uncertainty for this series of observations was 7 m s$^{-1}$. The excess RV scatter is consistent with the approximate 10 m s$^{-1}$ amplitude of solar-type oscillations derived from the scaling relations of Kjeldsen & Bedding (1995).

The atmospheric parameters of the stars under consideration were taken from P. Zielinski et al. (2009, in preparation, hereafter Z09), who have estimated their values using the method of Takeda et al. (2005a, 2005b). With these values and the luminosities estimated from available data, stellar masses were derived by fitting the ensemble of parameters characterizing the star ($\log(L/L_\odot)$, $\log T_{\text{eff}}$, $\log(g)$, and metallicity) to the evolutionary tracks of Girardi et al. (2000). The parameters of the three stars are summarized in Table 1.

Table 2

| Epoch (MJD) | RV (m s$^{-1}$) | $\sigma_{\text{RV}}$ (m s$^{-1}$) |
|------------|----------------|-------------------------------|
| 53546.31624 | 64.3           | 5.9                           |
| 53547.31606 | 50.8           | 7.4                           |
| 53551.30403 | 51.9           | 4.9                           |
| 53591.20234 | −30.9          | 9.3                           |
| 53593.19739 | −13.5          | 8.1                           |
| 53664.17558 | −6.3           | 4.2                           |
| 53868.43413 | −47.9          | 6.1                           |
| 53880.40891 | −38.9          | 5.3                           |
| 53893.36620 | −35.4          | 7.2                           |
| 53924.28339 | −20.4          | 9.0                           |
| 53934.43859 | −23.2          | 5.2                           |
| 53936.43783 | −27.9          | 7.7                           |
| 53955.20445 | 11.5           | 10.7                          |
| 53969.34923 | −10.2          | 8.8                           |
| 53996.26966 | 25.9           | 9.5                           |
| 54029.18670 | 49.8           | 6.8                           |
| 54041.14676 | 44.8           | 8.4                           |
| 54221.46062 | −11.8          | 6.4                           |
| 54248.39565 | 12.0           | 5.9                           |
| 54266.34430 | 7.8            | 10.0                          |
| 54318.39645 | 35.0           | 6.9                           |
| 54332.15993 | 66.5           | 5.9                           |
| 54346.12483 | 36.5           | 8.5                           |
| 54361.27718 | 24.7           | 8.5                           |
| 54371.24284 | −14.8          | 6.8                           |
| 54373.23936 | −29.4          | 11.1                          |
| 54373.25145 | −8.5           | 7.0                           |
| 54379.24422 | −26.4          | 8.3                           |
| 54381.21798 | −37.4          | 6.7                           |
| 54390.20274 | −41.7          | 6.6                           |
| 54587.45749 | 39.8           | 9.6                           |
| 54611.41703 | 28.7           | 6.9                           |
| 54617.37691 | 4.7            | 10.6                          |
| 54640.31703 | 3.4            | 10.7                          |
| 54665.26010 | −58.8          | 10.9                          |
| 54688.38071 | −72.2          | 9.7                           |
| 54703.35820 | −89.2          | 17.9                          |
| 54723.29658 | −59.6          | 11.4                          |
| 54739.24726 | −37.7          | 12.3                          |
| 54740.25007 | −56.6          | 7.1                           |
| 54756.19072 | −5.6           | 10.2                          |
| 54771.15219 | −27.1          | 9.3                           |
| 54811.05137 | 38.4           | 8.5                           |

2.1. BD+14 4559

RVs of BD+14 4559 were measured at 43 epochs over the period of 1265 days between MJ 53546 and 54811 (Table 2). Typically, the signal-to-noise ratio per resolution element in the spectra was 150–260 at 594 nm in 10–25 minutes of integration, depending on the atmospheric conditions. The estimated mean RV uncertainty for this star was 8 m s$^{-1}$.

BD+14 4559 (AG +14 2370, HIP 104780, LTT 16221) is a high proper motion K5 (Turon et al. 1993) star with $V = 9{\pm}63$ (Weis 1986), $B-V = 0{\pm}98 \pm 0.08$ (Høg et al. 2000), and $\pi = 19.99 \pm 1.61$ mas (Perryman & ESA 1997). The trigonometric parallax is consistent with the photometric value of $\pi = 24$ mas derived by Weis (1986). The atmospheric parameters of BD+14 4559 as determined in Z09 indicate that it is a dwarf with $\log(g) = 4.60 \pm 0.05$, $T_{\text{eff}} = 5008 \pm 20$ K, and [Fe/H] = 0.10 ± 0.07.

Using the existing Hipparcos (Perryman & ESA 1997) and Two Micron All Sky Survey (2MASS; Skrutskie et al. 2006) data, the substellar-mass companion to BD+14 4559 can be accurately determined.
Typically, the signal-to-noise ratio per resolution element in the period of 1655 days between MJD 53187 and 54842 (Table 3).

The derived absolute bolometric magnitude of the star, $M_V$, and log($L/L_\odot$), are 5.56 and $-0.32$, respectively. By comparing stellar parameters to evolutionary tracks of Girardi et al. (2000), we have computed the mass of BD+14 4559 to be $M/M_\odot = 0.86 \pm 0.15$ and $R/R_\odot = 0.95 \pm 0.2$.

The projected rotational velocity of BD+14 4559, $v \sin i = 2.5 \pm 1$ km s$^{-1}$, was estimated using the cross-correlation method (Benz & Mayor 1984). From this value and the adopted stellar radius we have obtained an estimate of the rotation period of $P_{\text{rot}} = 19$ days, which is much shorter than the observed 267 day period of the RV variations. From the uncertainty of the $v \sin(i)$ determination alone, the stellar rotation period may range from 13 to 32 days.

### 2.2. HD 240210

RVs of HD 240210 were measured at 38 epochs over the period of 1655 days between MJD 53187 and 54842 (Table 3). Typically, the signal-to-noise ratio per resolution element in the spectra was 150–300 at 594 nm in 4–10 minutes of integration, depending on the atmospheric conditions. The estimated mean RV uncertainty for this star was 8 m s$^{-1}$.

HD 240210 (BD+56 2959) is a K7 (Cannon & Pickering 1924) star with $V = 8^m33$, $B - V = 1^m\!65 \pm 0.04$ (Høg et al. 2000). The parallax of $\pi = 7.6 \pm 8.7$ mas (Perryman & ESA 1997) is consistent with the $\pi = 7.0 \pm 2.6$ mas of Gatwood et al. (1992). The atmospheric parameters of HD 240210 (Z09) indicate that the star is a giant with log($g$) = 2.31 $\pm$ 0.11, $T_{\text{eff}} = 4297 \pm 25$ K, and [Fe/H] = $-0.18 \pm 0.12$.

Using the existing Hipparcos (Perryman & ESA 1997) and 2MASS (Skrutskie et al. 2006) photometry, and the empirical calibrations of Ramírez & Meléndez (2005), we have estimated the effective temperature of the star to be $T_{\text{eff}} = 4290 \pm 94$ K. These values of effective temperature and log($g$) are consistent with a K3-giant star.

The derived absolute bolometric magnitude of the star and log($L/L_\odot$) are $0.86$ and $1.55$, respectively. By comparing stellar parameters to evolutionary tracks of Girardi et al. (2000) we have computed the mass of HD 240210 to be $M/M_\odot = 0.82 \pm 0.15$ and $R/R_\odot = 10.5 \pm 3$.

The projected rotational velocity of HD 240210, $v \sin i < 1.0 \pm 1$ km s$^{-1}$, was estimated using the cross-correlation method. From this value and the adopted stellar radius we have obtained an estimate of the rotation period of $P_{\text{rot}} > 530$ days.

### 2.3. BD+20 2457

RVs of BD+20 2457 were measured at 37 epochs over the period of 1833 days between MJD 53033 and 54866 (Table 4). Typically, the signal-to-noise ratio per resolution element in the spectra was 140–260 at 594 nm in 12–30 minutes of integration, depending on the atmospheric conditions. The estimated mean RV uncertainty for this star was 7 m s$^{-1}$.

BD+20 2457 (AG +20 1166) is a K2 (Heckmann 1975) star with $V = 9^m75$, $B - V = 1^m\!25 \pm 0.09$ (Høg et al. 2000) which is similar to $V = 9.57$ and $B - V = 1.35$ as measured photometrically by Busko et al. (1980), who observed it as a comparison star in their stellar variability study. The parallax of the star, $\pi = 5.0 \pm 26.50$ mas (Perryman & ESA 1997), is obviously very poorly determined. The atmospheric parameters of BD+20 2457 (Z09) indicate that it is a dwarf with log($g$) = 1.51 $\pm$ 0.05, $T_{\text{eff}} = 4137 \pm 10$ K, and [Fe/H] = $-1.00 \pm 0.07$.

Using the existing Hipparcos (Perryman & ESA 1997) and 2MASS (Skrutskie et al. 2006) photometry and the empirical calibrations of Ramírez & Meléndez (2005), we have estimated the effective temperature of the star to be $T_{\text{eff}} = 4127 \pm 94$ K. We conclude that the spectral type of BD+20 2457 is K2II.

The derived absolute bolometric magnitude of the star, $M_V$, and log($L/L_\odot$), are $-3.0$ and $3.17$, respectively. By comparing the stellar parameters to evolutionary tracks of Girardi et al. (2000), we have computed the mass of BD+20 2457 to fall in the $1.3$–$4.3$ $M/M_\odot$ range ($2.8 \pm 1.5$ $M/M_\odot$), and $R/R_\odot = 49 \pm 26$.

As for the other two stars, the projected rotational velocity of BD+20 2457, $v \sin i < 0.1$ km s$^{-1}$, was estimated using the cross-correlation method. This estimate is however very uncertain. We therefore assume as the upper limit of $v \sin i$ for BD+20 2457 the mean value of the rotational velocity for K2II stars of $v \sin i = 2.5$ km s$^{-1}$ from de Medeiros et al. (1996).

From this value and the adopted stellar radius we have obtained an estimate of the rotation period of the star, $P_{\text{rot}} > 983$ days.
3. ANALYSIS OF THE RV DATA

The process of modeling the RV measurements of the three stars discussed in this paper is documented in Figures 2–6. In the first step of the analysis, single, six-parameter, Keplerian orbits were least-squares fitted to data using the Levenberg–Marquardt algorithm (Press et al. 1992). Errors in the best-fit Keplerian parameters were estimated from both the parameter covariance matrix and bootstrap calculations involving 100 simulated data sets. For the K-giants, HD 240210 and BD+20 2457, the estimated, 7–8 m s$^{-1}$ errors in RV measurements were clearly insufficient to account for the observed RV variations in the post-fit residuals. This excess RV variability is believed to originate in fluctuations of the stellar surface and radial and non-radial oscillations of the giant stars (de Ridder et al. 2009) with the typical RV amplitudes around 20 m s$^{-1}$ (Hekker et al. 2006). To account for this effect, we have quadratically added 10 m s$^{-1}$ and 30 m s$^{-1}$ to the calculated RV errors for HD 240210 and BD+20 2457, respectively.

Evidently, in all the three cases, the best-fit residuals from single orbit models leave non-random trends that need to be modeled further. A statistical significance of these trends was assessed by calculating false alarm probabilities (FAP) for a null hypothesis that the trends can be adequately accounted for by a single Keplerian orbit and noise. The FAPs were calculated with

### Table 4

| Epoch (MJD) | RV (m s$^{-1}$) | $\sigma_{RV}$ (m s$^{-1}$) |
|------------|----------------|--------------------------|
| 53033.46852 | 253.5          | 5.2                      |
| 53360.35712 | −48.4          | 5.8                      |
| 53695.45657 | 434.5          | 3.8                      |
| 53751.49777 | 270.8          | 11.8                     |
| 53764.26867 | 224.4          | 5.2                      |
| 53774.23093 | 226.4          | 6.4                      |
| 53834.26977 | −206.2         | 4.4                      |
| 53867.17424 | −535.0         | 4.4                      |
| 53886.13781 | −431.2         | 7.2                      |
| 54044.49085 | 191.4          | 4.6                      |
| 54080.38513 | 224.7          | 7.3                      |
| 54105.31918 | 224.8          | 8.4                      |
| 54106.31976 | 232.0          | 7.2                      |
| 54109.31078 | 323.1          | 7.7                      |
| 54109.32586 | 330.4          | 10.1                     |
| 54130.24993 | 242.1          | 7.7                      |
| 54153.18142 | 105.1          | 8.3                      |
| 54160.16793 | 85.4           | 6.1                      |
| 54179.11821 | 100.4          | 4.7                      |
| 54192.28757 | 60.9           | 8.8                      |
| 54212.23429 | −63.9          | 5.7                      |
| 54242.15960 | −185.7         | 6.1                      |
| 54433.43142 | 271.7          | 7.5                      |
| 54464.33911 | 230.7          | 5.6                      |
| 54515.20532 | 64.8           | 4.9                      |
| 54570.25941 | −233.1         | 4.4                      |
| 54599.19278 | −264.9         | 11.6                     |
| 54602.16715 | −419.6         | 10.1                     |
| 54604.16879 | −303.6         | 6.9                      |
| 54774.49799 | 213.9          | 8.2                      |
| 54789.47046 | 261.9          | 7.1                      |
| 54790.45754 | 280.0          | 7.5                      |
| 54808.40943 | 535.4          | 9.5                      |
| 54838.33812 | 414.5          | 10.2                     |
| 54840.52386 | 341.7          | 7.6                      |
| 54842.31586 | 317.8          | 8.1                      |
| 54866.44103 | 299.2          | 7.2                      |

The process of modeling the RV measurements of the three stars discussed in this paper is documented in Figures 2–6. In the first step of the analysis, single, six-parameter, Keplerian orbits were least-squares fitted to data using the Levenberg–Marquardt algorithm (Press et al. 1992). Errors in the best-fit Keplerian parameters were estimated from both the parameter covariance matrix and bootstrap calculations involving 100 simulated data sets. For the K-giants, HD 240210 and BD+20 2457, the estimated, 7–8 m s$^{-1}$ errors in RV measurements were clearly insufficient to account for the observed RV variations in the post-fit residuals. This excess RV variability is believed to originate in fluctuations of the stellar surface and radial and non-radial oscillations of the giant stars (de Ridder et al. 2009) with the typical RV amplitudes around 20 m s$^{-1}$ (Hekker et al. 2006). To account for this effect, we have quadratically added 10 m s$^{-1}$ and 30 m s$^{-1}$ to the calculated RV errors for HD 240210 and BD+20 2457, respectively.

Evidently, in all the three cases, the best-fit residuals from single orbit models leave non-random trends that need to be modeled further. A statistical significance of these trends was assessed by calculating false alarm probabilities (FAP) for a null hypothesis that the trends can be adequately accounted for by a single Keplerian orbit and noise. The FAPs were calculated with the aid of the RV scrambling method (Wright et al. 2007, and references therein). As illustrated in Figure 6, FAP < 0.1% for the three data sets, which indicates that the additional, systematic RV variations seen in the data are statistically significant.

#### 3.1. BD+14 4559

As shown in Figure 2, the best-fit, single Keplerian orbit model for the RV variations of this star leaves behind a
significant, non-linear trend, which reaches the amplitude of \( \sim 50 \text{ m s}^{-1} \) and results in \( \sigma_{\text{RV}} = 17 \text{ m s}^{-1} \) over the 1300 day span of observations. At this point, rather than attempting to search the \( \chi^2 \)-space for a range of acceptable two-planet solutions, we have modeled the observed RV variations with a Keplerian orbit and a parabolic trend added to it. The best-fit \( \chi^2 \) value for this model drops from 4.4 for the single orbit case to 2.1 indicating, together with the FAP < 0.1%, that the trend is significant.

The final solution (Table 5) includes a planet with the minimum mass of \( 1.5 \text{ M}_{\text{Jup}} \), in a 269 day, \( 0.78 \text{ AU} \), \( e = 0.29 \) orbit around the star. Given the stellar mass of \( 0.86 \text{ M}_{\odot} \), the mass function of \( 2.5 \times 10^{-8} \text{ M}_{\odot} \), and the above RV amplitude and time span of the observed trend, one can broadly constrain a putative second companion to have a minimum mass of \( > 2.4 \text{ M}_{\text{Jup}} \) and the orbital radius \( > 2.3 \text{ AU} \), assuming \( e = 0 \).

The remaining post-fit rms residual, \( \sigma_{\text{RV}} = 11 \text{ m s}^{-1} \), slightly exceeds the estimated \( 8 \text{ m s}^{-1} \) precision of our RV measurements for this star. This excess “jitter” is within the estimated, intrinsically generated RV noise of \( 5 \text{ m s}^{-1} \) due to solar-type oscillations in stable dwarfs (Wright 2005).

### 3.2. HD 240210

A single orbit fit to the RV data for this star (Figure 3, Table 6) gives \( \chi^2 = 9.8 \), the post-fit rms residual, \( \sigma_{\text{RV}} = 39 \text{ m s}^{-1} \), and the very clear unmodeled RV variations. Our search for a
multiple orbit solution has resulted in several two- and three-planet models, which gave nearly identical improvements of the fit. This is illustrated in Figure 3 by a two-planet model, in which produces $\chi^2_r = 5.0$, $\sigma_RV = 25 \, \text{m s}^{-1}$, FAP $< 0.1 \%$, and the persistent lower-level systematics in the RV residuals. Because the existing measurements are obviously insufficient to obtain satisfactory constraints on the Keplerian model of the RV variations of this star, we list in Table 6 the initial one-planet solution and postpone further analysis until the time, when enough data become available. The provisional, best-fit orbit involves a 5.2 $M_{\text{Jup}}$ planet in a 502 day orbit with the semimajor axis of 1.16 AU and $e = 0.15$.

We also note that the 25 m s$^{-1}$ post-fit residual from the multiplanet fits is about 3 times larger than the formal precision of the RV measurements for HD 240210. As mentioned above, this is in agreement with the estimated RV scatter in K-giants due to their internal activity.

### 3.3. BD+20 2457

A single, Keplerian orbit fit to the RV data for this star is shown in Figure 4. The residuals from this fit are characterized by $\chi^2_r = 13.7$, $\sigma_RV = 105 \, \text{m s}^{-1}$, and a 600 day periodicity with the amplitude in excess of 100 m s$^{-1}$. Once again, as shown in Figure 6, the FAP for the existence of the second periodicity is less than 0.1%.

As shown in Figure 5, the best-fit, two-orbit solution offers a dramatic improvement to the Keplerian model of the RV variations observed in BD+20 2457, with the values of $\chi^2_r$ and $\sigma_RV$ reduced to 5.2 and 60 m s$^{-1}$, respectively. The model parameters calculated for the stellar mass of 2.8 $M_{\odot}$ and listed in Table 7 call for a system of two substellar-mass companions with the respective minimum masses of 21.4 $M_{\text{Jup}}$ and 12.5 $M_{\text{Jup}}$, orbital periods of 380 and 622 days, semimajor axes of 1.45 and 2.01 AU, and mild eccentricities of 0.15 and 0.18.

As in the case of HD 240210, the estimated, 7 m s$^{-1}$ uncertainty of the RV measurements of BD+20 2457 is much smaller than the actual, 60 m s$^{-1}$ rms of the post-fit residuals. This additional RV “jitter” broadly agrees with the 138 m s$^{-1}$ amplitude of the solar-type oscillations extrapolated for this star from Kjeldsen & Bedding (1995).

### 4. STELLAR PHOTOMETRY, ROTATION, AND LINE BISECTOR ANALYSIS

In order to verify that the observed RV periodicities are indeed caused by the Keplerian motion, we have thoroughly examined the existing photometry data in search for any periodic light variations, and performed a complete analysis of line bisectors and curvatures for each of the three stars discussed in this paper. This analysis used the cross-correlation method proposed by Martínez Fiorenzano (2005). For each star, the cross-correlation functions were computed from 1000 line profiles with the I$_2$ lines removed from the spectra. This procedure has been described in detail in Nowak et al. (2009).

Because both the Ca ii K emission line and the infrared Ca ii triplet lines at 849.8–854.2 nm are outside the range of our spectra, we have used the Hz line as a chromospheric activity indicator. As the shape of the Hz line in the spectra that showed significant telluric line contamination could be affected by it, care was taken to omit such spectra from the analysis. Equivalent widths (EW) of the line profile, defined as $I/I_c < 0.9$, were measured within the central 1 Å of the absorption profile, where the signature of stellar activity should be most pronounced (the uncertainty of the EW defined this way depends on the actual depth of the Hz line and the signal-to-noise ratio).

#### 4.1. BD+14 4559

The existing photometric databases for this star cover a wide range of epochs from MJD 47891 to 54792, and include extensive time series of measurements from the Hipparcos and Tycho (Perryman & ESA 1997), the NSVS (Woźniak et al. 2004), and the ASAS experiment (Pojmański 1997). None of the four data sets reveal any periodic light variations. In particular, the ASAS data, which are contemporaneous with our RV measurements of BD+14 4559, include observations made at 280 epochs between MJD 52754 and MJD 54792. In this case, the light variations are characterized by the mean value of $V = 9.650 \pm 0.021$ mag, which is consistent with the earlier measurements. As an example, the Lomb–Scargle (LS) periodogram of the ASAS photometry data is shown in Figure 7.

We have obtained 43 bisector and line curvature measurements, which resulted in the mean bisector velocity span, BVS = 4.6 ± 32 m s$^{-1}$, and the mean bisector curvature, $BC = -8 \pm 27$ m s$^{-1}$ with no trace of any periodic variability. No correlation between the BVS, BC, and RV was found ($r = 0.09$ and $r = 0.08$, respectively). Similarly, our analysis of 23 spectra of the star has shown that the EW of the Hz line exhibited a 3% rms scatter around its mean value of 396 ± 12 mÅ, which was not correlated with the observed RV variability ($r = -0.15$, while the critical value of the Pearson correlation coefficient at the confidence level of 0.01 is $r_{21,0.01} = 0.53$). The LS periodograms of the RVs, Hz EW, the BVS, and photometry for BD+14 4559 are shown in Figure 7. Clearly, the only observable that exhibits periodic time variations in time is the RV.

Using the scatter seen in the ASAS photometry of the star, and its rotational velocity determined above, we can estimate the amplitude of RV variations and of the BVS due to a possible presence of a spot on the stellar surface (Hatzes 2002). The
observed RV amplitude of BD+14 4559 is 3 times larger than the 35 m s\(^{-1}\) total RV amplitude predicted by Hatzes (2002) for a spot with a filling factor of \(f = 0.02\), on a star rotating at \(v \sin i = 2.5\) km s\(^{-1}\). Within the estimated uncertainty of the rotation velocity, the expected RV amplitude induced by a hypothetical spot is at least 2 times less than the observed one. The expected bisector variations of 5 m s\(^{-1}\) are comparable to the precision of our RV measurements and cannot have a detectable effect on our results. To summarize, the facts that there is no variability in the available photometric measurements over the period of \(\sim 19\) yr, the observed period of the RV variations is much longer than the estimated rotation period of the star, a significant eccentricity must be included in the best-fit model of the RV curve, and that the observed RV variations have been consistent over the five consecutive cycles of the period, make their interpretation in terms of a rotating spot very unlikely.

### 4.2. HD 240210

In the case of this star, the only long-term photometry available comes from 129 Tycho measurements with the \(V_T\) filter (Perryman & ESA 1997), between MJD 47946 and 49016. These data give the mean \(V_T = 8.499 \pm 0.012\) and the scatter in \(V_T\) of 0.136 mag. No periodic variability has been detected in these measurements. In addition, there are 19 epochs of photometric observations of this star available from the NSVS (Woźniak et al. 2004), collected between MJD 51362 and 51493. These data are characterized by the median \(V = 7.965 \pm 0.01\) and the scatter of 0.038. Since the NSVS photometry is of a much better quality than the Tycho data, we consider this scatter more believable. As shown in Figure 8, the LS periodogram of the Tycho photometry of HD 240210 reveals no significant periodicities.

The mean values of the BVS and the BC for this star are 62 ± 23 m s\(^{-1}\) and 23 ± 27 m s\(^{-1}\), respectively, with no statistically significant variability. No correlations between the RV residuals and the BVS and BC were found. The respective correlation coefficients are \(r = 0.08\) for the BVS and \(r = -0.18\) for the BC \((r_{36,0.01} = 0.40)\). Furthermore, the average EW of the H\(\alpha\) line measured over 28 epochs amounts to 602 ± 13, and the EW itself does not show any periodicities within 2%. We have considered a possible rotational modulation in H\(\alpha\) by calculating the Pearson correlation coefficient with the RV data residuals after removing the putative orbit. We found no correlation between the H\(\alpha\) EW and the RV residuals \((r = -0.26, \text{ whereas } r_{26,0.01} = 0.48)\). The LS periodograms of the four observables for HD 240210 discussed above are shown in Figure 8. As in the case of BD+14 4559, RV is the only one that varies periodically.

The observed photometric scatter in HD 240210, if interpreted in terms of a spot rotating with the star would also be detectable in both the RV and the line bisector data. The estimated total RV variations should be of the order of 23 m s\(^{-1}\) and the bisector variations should amount to 31 m s\(^{-1}\) (Hatzes 2002). The semi-amplitude of the observed periodic signal in RV is much larger than the predicted variations induced by a hypothetical spot of size consistent with the existing photometric variations. Periodic variations in line bisectors produced by a hypothetical spot should manifest themselves as the uncertainty in our BVS measurements, which is comparable to the predicted value. In any case, these variations should show a periodicity equal to the stellar rotation period, \(P > 530\) days. Because no
such variations are present in our bisector measurements for this star, we conclude that the observed scatter is a combination of low quality photometry and possible solar-type oscillations of the star. Such oscillations might account for 3.3 mmag in the photometric variability and $V_{\text{osc}} = 13$ m s$^{-1}$ in the BVS using the scaling relations of Kjeldsen & Bedding (1995).

4.3. BD+20 2457

As in the case of BD+14 4559, we have identified four photometric data sets that we could utilize in our analysis, but we have concentrated on the data from the ASAS project (Pojmański 1997), which contains extensive measurements that are contemporaneous with our RV observations. These data include 120 measurements made between MJD 52622 and MJD 53400 and the 191 additional ones made between MJD 52623 and MJD 54911. We have used 319 grade A measurements from the ASAS to derive the mean $V = 9.693 \pm 0.051$ mag and to compute their LS periodogram in search for any periodicities. As shown in Figure 9, no significant periodicities are present in the spectrum. Similarly, no periodicities were found in the other three, lower quality time series.

We have obtained 37 bisector and curvature measurements for this star, with the mean BVS $= 105 \pm 39$ m s$^{-1}$, and the mean BC $= 47 \pm 30$ m s$^{-1}$, both revealing no significant periodicities in the corresponding LS periodograms (Figure 9). We have also searched for correlations between the RV residuals and the BVS and BC after having removed the putative companions b and c. The respective correlation coefficients are $r = -0.33$ and $r = -0.19$ for BVS and $r = -0.12$ and $r = -0.09$ for BC, respectively (the critical value of Pearson correlation coefficient at the confidence level of 0.01 is $r_{35,0.01} = 0.42$).

The EW of the Hα line for BD+20 2457 is $653 \pm 23$, and this parameter shows no significant periodicities within 3.5% for 23 observations (Figure 9). As for the other two stars, we have searched for a possible rotational modulation in Hα by calculating the Pearson correlations coefficient with the RV data residuals after removing the individual orbits. We found that there is no significant correlation between the Hα EW and the RV residuals ($r = -0.21$ and $-0.51$, respectively, and $r_{21,0.01} = 0.53$).

A photometric scatter resulting from a spot rotating with BD+20 2457 would induce the estimated total RV variations on the order of $84$ m s$^{-1}$ and the bisector variations of $11$ m s$^{-1}$ (Hatzes 2002). The semi-amplitude of the observed periodic signal in RVs is much larger. As in BD+14 4559, the anticipated variability in line bisectors produced by a periodic spot rotation is comparable to the precision of our RV measurements and cannot have a detectable effect on our results. These variations, if real, should show a periodicity equal to the star’s rotation period, $P > 983$ days, which appears not to be the case. Since no measurable variations are present in the bisectors and line curvatures, we conclude that the observed photometric scatter is again a combination of low quality photometry and possible solar-type oscillations of the star. Such oscillations might account for 5.9 mmag and up to $V_{\text{osc}} = 138$ m s$^{-1}$ using the scaling relations of Kjeldsen & Bedding (1995).

5. DISCUSSION

The most recent results from the PTPS search for extrasolar planets described in this paper add a Jupiter-mass planet to the large body of planets around main-sequence stars (Santos 2008) and increase the still modest number of substellar-mass objects around GK-giants (e.g., Döllinger et al. 2009; Liu et al. 2009; Niedzielski et al. 2007, 2009; Sato et al. 2008b, and references therein) by at least three new detections.

In principle, the long period RV variations in red giants may be related to non-Keplerian effects, including the stellar rotation modulation of surface inhomogeneities, intrinsic activity, or non-radial pulsations (e.g., Hatzes et al. 2006). Our analysis of the available photometric data and the behavior of line bisectors in the three stars, following the established practices (Queloz et al. 2001), has shown no significant correlations between the time variability of these stellar activity indicators and the RV variations. Consequently, our RV measurements find the most plausible explanation in terms of the Keplerian motion of substellar-mass companions around the observed stars. The details of this analysis can be found in Niedzielski et al. (2008).

In the case of the K2-dwarf, BD+14 4559, the RV data reveal a planet with the minimum mass of $1.5 M_J$ in a 269 day, 0.78 AU, $e = 0.29$ orbit around the star. An additional, long-term trend seen in the data suggests the presence of another, long-period companion, but the time span of the available RV measurements is not yet sufficient to constrain its maximum mass with any confidence. If the third body turns out to be a planet, it will add to a growing number of confirmed and anticipated multiplanet systems around solar-type stars (Wright et al. 2009).

A single planet model of the RV variations in the K3-giant, HD 240210, leaves highly correlated post-fit residuals signaling a possible presence of additional substellar-mass bodies orbiting the star. The presently available data are not sufficient to obtain an unambiguous multibody solution for this system. The provisional parameters for one planet that can be fitted for give a $5.2 M_J$ body in a 502 day, 1.16 AU, $e = 0.15$ orbit that will...
The solid lines delimit the masses of the inner companion computed from the mass function over the estimated 1.3–4.3 $M_\odot$ range for the stellar mass and orbital inclinations from 18° to 90°. The dashed lines have the same meaning for the outer companion. The dotted lines constrain the theoretical mass range for brown dwarfs.

have to be revised, when another planet (or planets) are added to the current model. Nevertheless, the solution is robust enough to conclude that HD 240210 has at least one planet with large minimum mass. This result further emphasizes the observed correlation between the stellar mass and the masses of orbiting planets (Johnson et al. 2007; Lovis & Mayor 2007).

The RV variations observed in the K2-giant, BD+20 2457, can be unambiguously modeled by two Keplerian orbits with the periods of 380 and 622 days, 1.4 and 2 AU semimajor axes, and the respective eccentricities of 0.15 and 0.18. For the estimated stellar mass of 2.8 $M_\odot$, these parameters yield minimum masses of the orbiting bodies of 12.5 $M_J$ and 21.4 $M_J$, respectively. The ~1.6 ratio of the orbital periods places the orbits relatively close to the 3:2 mean motion resonance (MMR). These characteristics suggest that the two bodies share a common origin from a disk that once surrounded the primary star. At the same time, at least one of its companions, and very possibly both of them, have masses that exceed the 13 $M_J$ deuterium burning limit that conventionally separates planets from brown dwarfs.

Further analysis of this intriguing system is complicated by the fact that, due to the highly uncertain parallax of the primary, its mass is poorly known and spans a range that is at least as wide as 1.3–4.3 $M_\odot$ (Section 2). This is illustrated in Figure 10, which shows the constraints on the masses of the two substellar companions under the usual assumption of a random distribution of orbital inclinations. There is a 95% probability that the masses of the two bodies fall in the respective ranges delimited by the inclinations of 18° and 90°. Clearly, high masses, approaching the hydrogen burning limit, are very unlikely. Similarly, a possibility that both masses stay below the deuterium burning limit is marginal at best, unless one accepts an unrealistically low primary mass, approaching 1 $M_\odot$ or even less. Consequently, within the currently available constraints, it is reasonable to assume that the true masses of the two companions to BD+20 24 57 have masses in the brown dwarf range. An extensive study of the dynamics of this system will be described in a subsequent paper.

The two brown dwarf–mass bodies discovered by our survey around BD+20 24 57 add to the previous three detections of single objects of this kind around intermediate-mass stars (Hatzes et al. 2005; Lovis & Mayor 2007; Liu et al. 2008). Only one such two-companion system has been previously identified around a solar-mass star, HD 168443, by Marcy et al. (2001). On the basis of the facts presented above, it appears entirely reasonable to assume that the two companions to this star have originated from a massive disk and acquired enough mass to place them above the formal 13 $M_J$ brown dwarf limit. This raises an intriguing possibility that, in the case of substellar-mass companions to giants, both the concept of the brown dwarf desert (Marcy & Butler 2000) and the meaning of the deuterium burning limit will have to be revisited. In fact, this problem has been remarked on by Lovis & Mayor (2007) in the context of the observed correlation between the stellar and the planetary mass.

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