## Observation of $B \to \phi K$ Decays

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Abstract

We report the observation of the decay $B^\pm \to \phi\phi K^\pm$ and find evidence for $B^0 \to \phi\phi K^0$. These results are based on a 414 fb$^{-1}$ data sample that contains $449 \times 10^6 \ B\bar{B}$ pairs, collected with the Belle detector at the KEKB asymmetric-energy $e^+e^-$ (3.5 on 8 GeV) collider operating at the $\Upsilon(4S)$ resonance. This is the first observation of a $b \to s\bar{s}s\bar{s}$ transition. The branching fractions for these decay modes are measured to be $B(B^\pm \to \phi\phi K^\pm) = (3.2^{+0.6}_{-0.5} \pm 0.3) \times 10^{-6}$ and $B(B^0 \to \phi\phi K^0) = (2.3^{+1.0}_{-0.7} \pm 0.2) \times 10^{-6}$ for $\phi\phi$ invariant mass below 2.85 GeV/c$^2$. The corresponding partial rate asymmetry for the charged $B$ mode is measured to be $A_{CP}(B^\pm \to \phi\phi K^\pm) = 0.01^{+0.19}_{-0.16} \pm 0.02$. Results for other related charmonium decay modes are also reported. We also search for $CP$ asymmetry using the $\phi\phi$ candidates within the $\eta_c$ mass region. The value obtained is $A_{CP}(B^\pm \to \phi\phi K^\pm, M_{\phi\phi} \approx M_{\eta_c}) = 0.15^{+0.16}_{-0.17} \pm 0.02$, which is consistent with no asymmetry.

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Evidence of charmless \( B \to \phi\phi K \) decays has been reported by the Belle collaboration using \( 85 \times 10^6 B\bar{B} \) pairs \[^1\]. In the standard model (SM), this decay channel requires the creation of an additional final \( s\bar{s} \) quark pair in a \( b \to s\bar{s}s \) process, such as \( B \to \phi K \). Although the statistical error of the previous \( B \to \phi\phi K \) measurement was large, the branching fraction was found to be around 1/3 of that of \( B \to \phi K \), which provided useful information for understanding quark fragmentation in \( B \) decay. Moreover, it indicated that with sufficient statistics the decay \( B \to \phi\phi K \) could be used to search for a possible non-SM \( CP \)-violating phase in the \( b \to s \) transition \[^2\]. Direct \( CP \) violation could be enhanced as high as the 40\% level if there is sizable interference between the transitions due to non-SM physics and the decays via the \( \eta_c \) resonance.

In this paper, we report the observation of charmless and charmful \( B \to \phi\phi K \) decays based on a 414 fb\(^{-1} \) data sample that contains \( 449 \times 10^6 B\bar{B} \) pairs, collected with the Belle detector at the KEKB asymmetric-energy \( e^+e^- \) (3.5 on 8 GeV) collider \[^3\] operating at the \( \Upsilon(4S) \) resonance.

The Belle detector is a large-solid-angle magnetic spectrometer that consists of a silicon vertex detector (SVD), a 50-layer central drift chamber (CDC), an array of aerogel threshold Čerenkov counters (ACC), a barrel-like arrangement of time-of-flight scintillation counters (TOF), and an electromagnetic calorimeter comprised of CsI(Tl) crystals (ECL) located inside a super-conducting solenoid coil that provides a 1.5 T magnetic field. An iron flux-return located outside of the coil is instrumented to detect \( K^0_L \) mesons and to identify muons (KLM). The detector is described in detail elsewhere \[^4\]. Two inner detector configurations were used. A 2.0 cm radius beampipe and a 3-layer silicon vertex detector (SVD1) were used for the first sample of \( 152 \times 10^6 B\bar{B} \) pairs, while a 1.5 cm radius beampipe, a 4-layer silicon detector (SVD2) and a small-cell inner drift chamber were used to record the remaining 297 \( \times 10^6 B\bar{B} \) pairs \[^5\].

Charged kaons are required to have impact parameters within ±2 cm of the interaction point (IP) along the \( z \)-axis (anti-parallel to the positron direction) and within 0.2 cm in the transverse plane. Each track is identified as a kaon or a pion according to a \( K/\pi \) likelihood ratio, \( R(K/\pi) = \mathcal{L}_K/(\mathcal{L}_K + \mathcal{L}_\pi) \), where \( \mathcal{L}_K/\mathcal{L}_\pi \) is the likelihood of kaons/pions derived from the responses of TOF and ACC systems and the energy loss measurements from the CDC. The likelihood ratio is required to exceed 0.6 for kaon candidates; within the momentum range of interest, this requirement is 88\% efficient for kaons and has a misidentification rate for pions of 8.5\%. Neutral kaons are reconstructed via the decay \( K^0_S \to \pi^+\pi^- \) and have an invariant mass \( 0.482 \text{ GeV}/c^2 < M_{\pi^+\pi^-} < 0.514 \text{ GeV}/c^2 \) (±4\( \sigma \) mass resolution). The \( \pi^+\pi^- \) vertex is required to be displaced from the IP and the flight direction must be consistent with a \( K^0_S \) that originated from the IP. The required displacement increases with the momentum of the \( K^0_S \) candidate.

\( B \) meson candidates are reconstructed in the five kaon final state. Two kinematic variables are used to distinguish signal candidates from backgrounds: the beam-energy constrained mass \( M_{bc} = \sqrt{E_{\text{beam}}^2 - |\vec{P}_{\text{recon}}|^2} \) and the energy difference \( \Delta E = E_{\text{recon}} - E_{\text{beam}} \), where \( E_{\text{beam}} \) is the beam energy, and \( E_{\text{recon}} \) and \( \vec{P}_{\text{recon}} \) are the reconstructed energy and momentum of the signal candidate in the \( \Upsilon(4S) \) rest frame. The resolution of \( M_{bc} \) is approximately 2.8 MeV$/c^2$, dominated by the beam energy spread, while the \( \Delta E \) resolution is around 10 MeV. Candidates with five kaons within the region \( |\Delta E| < 0.2 \text{ GeV} \) and \( 5.2 \text{ GeV}/c^2 < M_{bc} < 5.29 \text{ GeV}/c^2 \) are selected for further consideration. The signal region is defined as \( 5.27 \text{ GeV}/c^2 < M_{bc} < 5.29 \text{ GeV}/c^2 \).

The dominant backgrounds are \( e^+e^- \to q\bar{q} \) (\( q = u, d, s, c \)) continuum events. Event
topology and $B$ flavor tagging are used to distinguish the jet-like continuum events and the spherically distributed $BB$ events. Seven event-shape variables are combined into a single Fisher discriminant $F$. The Fisher variables include the angle between the thrust axis of the $B$ candidate and the thrust axis of the rest of the event ($\cos \theta_T$), five modified Fox-Wolfram moments $R$, and a measure of the momentum transverse to the event thrust axis ($S_L$) $F$. Two other variables that are uncorrelated with the Fisher discriminant and help to distinguish signal from the continuum are $\cos \theta_B$, where $\theta_B$ is the angle between the $B$ flight direction and the beam direction in the $e^+e^-$ center of mass frame, and $\Delta z$, the $z$-vertex difference between the signal $B$ candidate and its accompanying $B$. We form signal and background probability density functions (PDFs) for the Fisher discriminant, $\cos \theta_B$ and $\Delta z$ using the signal Monte Carlo (MC) events and sideband data (5.2 GeV/$c^2 < M_{bc} < 5.26$ GeV/$c^2$), respectively. The products of the PDFs for these variables give signal and background likelihoods $L_S$ and $L_{BG}$ for each candidate, allowing a selection to be applied to the likelihood ratio $R = L_S/(L_S + L_{BG})$.

Additional background discrimination is provided by the quality of the $B$ flavor tagging of the accompanying $B$ meson. The standard Belle flavor tagging package gives two outputs: a discrete variable indicating the flavor of the tagging $B$ and dilution factor $r$, which ranges from zero for no flavor information to unity for unambiguous flavor assignment. The continuum background is reduced by applying a selection requirement of $R$ for events in each $r$ region according to the figure of merit defined as $N_S/\sqrt{N_S + N_{BG}}$, where $N_S$ denotes the expected $\phi\phi K$ signal yields based on MC simulation and the branching fraction reported in our previous measurements, and $N_{BG}$ denotes the expected $q\bar{q}$ yields from sideband data. This requirement removes (61-81)% of the continuum background while retaining (80-92)% of the signal, and depends on the decay channel ($\phi\phi K$ or $\phi\phi K^0$) and the SVD configuration during the measurement (SVD1 or SVD2). Backgrounds from other $B$ decays are investigated using a large MC sample and are found to be negligible after the $R$ requirement.

The signal yields are extracted by applying an unbinned extended maximum likelihood (ML) fit to the events with $M_{bc} > 5.2$ GeV/$c^2$ and $|\Delta E| < 0.2$ GeV. For the $\phi\phi K^\pm$ mode, we simultaneously obtain the yield and the partial rate asymmetry $A_{CP}$ using the likelihood, defined as:

$$L = \exp\left[-(N_S + N_B)\right] \prod_i \left(\sum_j \frac{1}{2} [1 - q_i \cdot \mathcal{A}_{ij}] N_j P_{ij}(M_{bc}, \Delta E)\right),$$

(1)

where $i$ is the identifier of the $i$-th event, $j$ indicates signal or background, $P(M_{bc}, \Delta E)$ is the two-dimensional PDF of $M_{bc}$ and $\Delta E$, and $q$ indicates the $B$ meson flavor, $+1$ for $B^+$ and $-1$ for $B^-$, respectively. For neutral $B$ events, $1/2[1 - q_i \cdot \mathcal{A}_{CP}]$ in Eq. (1) is replaced by 1.

The $M_{bc}$ PDFs are modeled by a Gaussian function for signals and an ARGUS function for the continuum, while a Gaussian is used to describe the signal $\Delta E$ and a 2nd order Chebyshev polynomial is used for the background $\Delta E$. The parameters of the PDFs are determined using high-statistics MC samples and sideband data for signal and background shapes. The signal PDFs are calibrated by comparing the $M_{bc}$ and $\Delta E$ distributions of the $B^+ \rightarrow D^0(K^+\pi^-\pi^-\pi^+)\pi^+$ events with the MC expectation.

We search for charmless $B \rightarrow \phi\phi K$ decays by requiring the $\phi\phi$ invariant mass ($M_{\phi\phi}$) to be less than 2.85 GeV/$c^2$, the region below the charm threshold. Candidate $\phi$ mesons are identified by requiring the invariant masses of $K^+K^-$ pairs ($M_{K^+K^-}$) to be in the range 1.0 GeV/$c^2$ to 1.04 GeV/$c^2$ ($\pm 4.6\sigma$). Figure shows the $M_{bc}$ and $\Delta E$ projections with the
fit curves superimposed. Clear signals appear in both $B^\pm$ and $B^0$ modes with signal yields of $37.0^{+6.9}_{-6.0}$ and $7.8^{+3.2}_{-2.5}$, respectively. Although $K^+K^-$ candidates are required to lie in the $\phi$ mass region, non-$\phi$ backgrounds may also contribute. Figure 2(a) shows the $M_{K^+K^-}$ vs. $M_{K^+K^-}$ distributions for $(K^+K^-K^+K^-)K^\pm$ candidates in the signal region, where the two $K^+K^-$ pairs are required to have invariant masses less than $1.2 \text{ GeV}/c^2$. Events in the two $\phi$ bands are used to estimate the $B^\pm \rightarrow \phi K^+K^-K^\pm$ contribution. Figure 2(b) shows $B$ signal yields $[[1]]$ as a function of the $K^+K^-$ invariant mass after requiring the other $K^+K^-$ pair to have a mass in the $\phi$ mass region. The $B$ signal yields are fitted with a threshold function in the region $0.98 \text{ GeV}/c^2 < M_{K^+K^-} < 1.2 \text{ GeV}/c^2$, excluding $\phi$ mass region ($1.0 \text{ GeV}/c^2 < M_{K^+K^-} < 1.04 \text{ GeV}/c^2$). The amount of the non-$\phi$ contribution is estimated by interpolating the $B$ yields in the $\phi$ sideband region to the $\phi$ mass region, which is $4.4^{+0.8}_{-0.7}$ events. Since events in the two $\phi$ bands contain both true $\phi$ mesons and non-resonant $K^+K^-$ pairs, the area underneath the $\phi$ mass region in Fig. 2(b) includes the $\phi K^+K^-K^\pm$ contribution but counts the non-resonant $5K$ component twice. Therefore, we estimate the non-resonant $B \rightarrow 5K$ contribution using the $B$ signal yield in the upper right corner of the dashed region in Fig. 2(a). We assume a phase-space distribution in 4-kaon mass. We obtain $1.3 \pm 0.4$ non-resonant events in the $\phi\phi K^\pm$ sample. After subtracting this contribution of $1.3 \pm 0.4$ events, the non-$\phi\phi K$ fraction is calculated to be $(7 \pm 4)\%$. The same procedure is applied to the $\phi\phi K^0$ sample; here we obtain a fraction of $(7 \pm 9)\%$.

Table I summarizes the $\phi\phi K$ results after subtracting the non-$\phi\phi K$ contribution. Signal efficiencies are obtained by generating $\phi\phi K$ MC events, where the same $M_{\phi\phi} < 2.85 \text{ GeV}/c^2$ requirement is applied. Systematic uncertainties in the fit are obtained by performing fits in which the signal peak positions and resolutions of the signal PDFs are successively varied by $\pm 1\sigma$. The quadratic sum of each deviation from the central value of the fit gives the systematic uncertainty of the fit. For each systematic check, the statistical significance is taken as $\sqrt{-2\ln(L_{\text{feeddown}}/L_{\text{max}})}$, where $L_{\text{feeddown}}$ and $L_{\text{max}}$ are the likelihoods at the expected non-$\phi\phi K$ yields and the best fit, respectively. We regard the smallest value as our significance including the systematic uncertainty. The number of $B^+B^-$ and $B^0\bar{B}^0$ pairs are assumed to be equal.

| Mode | Yields | efficiencies(%) | $\Sigma$ | $B(10^{-6})$ |
|------|--------|----------------|---------|-------------|
| $B^\pm \rightarrow \phi\phi K^\pm$ ($M_{\phi\phi} < 2.85 \text{ GeV}/c^2$) | $34.2^{+6.4}_{-5.8}$ | 2.41 | 9.5 | $3.2^{+0.6}_{-0.5} \pm 0.3$ |
| $B^0 \rightarrow \phi\phi K^0$ ($M_{\phi\phi} < 2.85 \text{ GeV}/c^2$) | $7.3^{+3.0}_{-2.4}$ | 0.69 | 4.7 | $2.3^{+1.0}_{-0.7} \pm 0.2$ |
| $B^\pm \rightarrow \eta_c K^\pm$, $\eta_c \rightarrow \phi\phi$ | $29.7^{+6.0}_{-5.5}$ | 2.72 | 7.2 | $2.4^{+0.6}_{-0.5} \pm 0.2$ |
| $B^\pm \rightarrow \eta_c K^\pm$, $\eta_c \rightarrow K^+K^-$ | $76.8^{+13.6}_{-12.4}$ | 4.85 | 9.4 | $3.5 \pm 0.6 \pm 0.3$ |
| $B^\pm \rightarrow \eta_c K^\pm$, $\eta_c \rightarrow 2(K^+K^-)$ | $104.6^{+20.2}_{-17.3}$ | 9.93 | 10.2 | $2.4^{+0.5}_{-0.4} \pm 0.2$ |
| $B^\pm \rightarrow J/\psi K^\pm$, $J/\psi \rightarrow \phi K^+K^-$ | $25.5^{+7.0}_{-6.0}$ | 4.67 | 8.5 | $1.2 \pm 0.3 \pm 0.1$ |
| $B^\pm \rightarrow J/\psi K^\pm$, $J/\psi \rightarrow 2(K^+K^-)$ | $41.0^{+7.3}_{-6.6}$ | 9.41 | 9.7 | $0.97^{+0.17}_{-0.16} \pm 0.1$ |

The performance of the $R$ requirement is studied by checking the data-MC efficiency ratio using the $B^+ \rightarrow \bar{D}^0(\rightarrow K^+\pi^-\pi^-\pi^+)\pi^+$ sample. The obtained error is 2.7-2.8%. The systematic errors on the charged track reconstruction are estimated to be around $1\%$ per track using partially reconstructed $D^*$ events. Therefore, the tracking systematic error is
FIG. 1: Distributions of $M_{bc}$ and $\Delta E$ for the decay modes $B^\pm \rightarrow \phi\phi K^\pm$ (a,b) and $B^0 \rightarrow \phi\phi K^0$ (c,d), with $\phi\phi$ invariant mass less then 2.85 GeV/$c^2$. The open histograms represent the data, the solid blue curves show the result of the fit, the dash-dotted red lines represent the signal contributions and the dashed green curves show the continuum background contributions.

5% (5 tracks) for the $\phi\phi K^+$ mode and 4% for the $\phi\phi K^0$ mode. The kaon identification efficiency is studied using samples of inclusive $D^+ \rightarrow D^0\pi^+, D^0 \rightarrow K^-\pi^+$ decays. The $K^0_S$ reconstruction is verified by comparing the ratio of $D^+ \rightarrow K^0_S\pi^+$ and $D^+ \rightarrow K^-\pi^+\pi^+$ yields. The resulting $K^0_S$ detection systematic error is 4.9%. The uncertainty in the number of $B\bar{B}$ events is 1%. The final systematic error is obtained by summing all correlated errors linearly and then quadratically summing the uncorrelated errors.

After subtracting the non-$\phi\phi K$ contribution, the branching fractions for charmless $B \rightarrow \phi\phi K$ decays are $B(B^\pm \rightarrow \phi\phi K^\pm) = (3.2^{+0.6}_{-0.5} \pm 0.3) \times 10^{-6}$ with 9.5$\sigma$ significance and $B(B^0 \rightarrow \phi\phi K^0) = (2.3^{+1.0}_{-0.7} \pm 0.2) \times 10^{-6}$ with 4.7$\sigma$ significance. The measured charge asymmetry for $B^\pm \rightarrow \phi\phi K^\pm$ decay is $0.01^{+0.19}_{-0.16} \pm 0.02$. The first error is statistical and the second is systematic.

It is of interest to search for possible $\phi\phi$ resonances above charm threshold. Figure 3(a) shows the $B$ signal yields divided by the bin size as a function of $M_{\phi\phi}$ after releasing the $M_{\phi\phi} < 2.85$ GeV/$c^2$ requirement. There is no specific enhancement in high $\phi\phi$ mass region except for the $\eta_c$ peak near 3 GeV/$c^2$. Reference 2 suggests a large $CP$ asymmetry, that arises from the interference between $B^\pm \rightarrow \phi\phi K^\pm$ and $B^\pm \rightarrow \eta_c(\rightarrow \phi\phi)K^\pm$ decays. The
FIG. 2: (a) The distribution of $M_{K^+K^-}$ vs. $M_{K^+K^-}$ for the $K^+K^-K^+K^-K^\pm$ candidates in the $M_{bc} - \Delta E$ signal box with $M_{K^+K^-} < 1.2$ GeV/$c^2$. The two $K^+K^-$ bands indicate the $\phi$ mass region ($1.0$ GeV/$c^2 < M_{K^+K^-} < 1.04$ GeV/$c^2$). The horizontal and vertical dashed lines are located at $M_{K^+K^-} = 1.05$ GeV/$c^2$. The rectangle on the upper right is the $\phi\phi$ sideband region; events in this region are used to estimate the non-resonant $B \to 5K$ contribution. (b) $B$ signal yields as a function of the $M_{K^+K^-}$ after requiring the other $K^+K^-$ pair has a mass in the $\phi$ mass region. The threshold function is used to be the fit curve and events with $1.0$ GeV/$c^2 < M_{K^+K^-} < 1.04$ GeV/$c^2$ are excluded when the fit is performed.

events with $\phi\phi$ invariant mass within $\pm 40$ MeV/$c^2$ of the nominal $\eta_c$ mass are selected to investigate this asymmetry. The measured $CP$ asymmetry is $0.15^{+0.10}_{-0.17} \pm 0.02$, which is consistent with no asymmetry.

We study possible charmonium states by checking $B$ yields with $M_{4K}$ between $2.8$ GeV/$c^2$ and $3.2$ GeV/$c^2$. Since $\eta_c$ and $J/\psi$ mesons may decay to $\phi K^+K^-$ and $2(K^+K^-)$ pairs, a mass scan is performed with and without the requirement that the $K^+K^-$ pair lie in the $\phi$ mass region. As shown in Fig. 3 clear $\eta_c$ and $J/\psi$ resonances are visible in the $\phi K^+K^-$ and $4K$ samples while only an $\eta_c$ peak appears in the $\phi\phi$ mode.

We obtain the signal yields for $B^\pm \to \eta_c K^\pm$ and $B^\pm \to J/\psi K^\pm$ by performing binned histogram fits to Figs. 3(b), 3(c) and 3(d). The $J/\psi$ signal PDF is modeled with a Gaussian function while the $\eta_c$ PDF is described by a Breit-Wigner function convolved with a Gaussian resolution function, which has the same Gaussian width as the $J/\psi$ PDF. Since sizable signals are observed in the $4K$ mode, the parameters are determined using the $4K$ sample and the same signal PDFs are then applied to the $\phi K^+K^-$ and $\phi\phi$ samples. The obtained Gaussian width is measured to be $4.0^{+1.0}_{-0.8}$ MeV/$c^2$ by performing a fit with a $2^{nd}$ order Chebyshev polynomial as the non-resonant PDF after excluding events in the $\eta_c$ mass region ($2.94$ GeV/$c^2 < M_{4K} < 3.02$ GeV/$c^2$). A fit is performed to the full range with the Gaussian width fixed. The signal yields are summarized in Table III. The peak positions obtained for the $\eta_c$ and $J/\psi$ are $2.979 \pm 0.002$ GeV/$c^2$ and $3.094 \pm 0.001$ GeV/$c^2$, respectively, consistent with the nominal $\eta_c$ and $J/\psi$ masses. The $\eta_c$ Breit-Wigner width is measured to be $25.2^{+7.7}_{-6.0} \pm 0.3$ MeV/$c^2$, where the central value is consistent with the world average and the second error is due to the uncertainty in the mass resolution.

For the $\phi K^+K^-$ and $\phi\phi$ modes, the non-$\phi$ contribution is determined from the $B$ signal yields for events with one $K^+K^-$ pair in the $\phi$ sideband region ($1.05$ GeV/$c^2 < M_{K^+K^-} <
FIG. 3: $B$ signal yields as a function of (a,b) $M_{\phi\phi}$, (c) $M_{\phi K^+ K^-}$ and (d) $M_{4K}$. In (a) we use different bin sizes for $M_{\phi\phi}$ less than 3 GeV/$c^2$ and greater than 3 GeV/$c^2$. The subset with $M_{\phi\phi}$ from 2.8 GeV/$c^2$ to 3.2 GeV/$c^2$ is shown in (b).

1.09 GeV/$c^2$ and the $4K$ and $\phi K^+ K^-$ masses are in the charmonium resonance region, respectively. We find $3.0^{+2.0}_{-1.4}$ events in the $\eta_c \to \phi\phi$ mode, $6.4^{+5.4}_{-4.4}$ events in the $\eta_c \to \phi K^+ K^-$ mode, and $3.4^{+3.6}_{-2.6}$ in the $J/\psi \to \phi K^+ K^-$ mode. After subtracting the feed-down yields, we obtain the results listed in Table I.

Signal efficiencies are determined using signal MC and the detection systematic uncertainties are similar to what was described in the charmless $\phi\phi K$ part. Fit systematic uncertainties are estimated by successively varying the peak positions and resolutions of the $M_{bc} - \Delta E$ signal PDFs as well as the convolution Gaussian width in the fit. The quadratic sum of each deviation gives the fit systematic errors. Since the sub-decay branching fractions of $\eta_c$ and $J/\psi$ mesons to $4K$, $\phi K\bar{K}$ and $\phi\phi$ final states are not precisely known, we provide the product of branching fractions for various decays in Table I. Using the known branching fractions of $\mathcal{B}(B^\pm \to \eta_c K^\pm) = (9.1 \pm 1.3) \times 10^{-4}$ and $\mathcal{B}(B^\pm \to J/\psi K^\pm) = (1.008 \pm 0.035) \times 10^{-3}$, the subdecay branching fractions are calculated and listed in Table I.

In summary, we have observed the charmless decay $B^\pm \to \phi\phi K^\pm$ and evidence of $B^0 \to \phi\phi K^0$. We also report the $CP$ asymmetry of the charged decay and measurements of other closely related charmonium decays. The results are consistent with the previous measurements, but have considerably improved precision due to the increase in statistics.

9
TABLE II: The measured branching fractions of secondary charmonium decays and the world averages [12].

| Decay mode       | $\mathcal{B}$ (measured)      | $\mathcal{B}$ (PDG 2006 value) |
|------------------|--------------------------------|---------------------------------|
| $\eta_c \to \phi\phi$ | $(2.7^{+0.6}_{-0.5} \pm 0.4) \times 10^{-3}$ | $(2.7 \pm 0.9) \times 10^{-3}$ |
| $\eta_c \to \phi K^+K^-$ | $(3.9^{+0.7}_{-0.6} \pm 0.6) \times 10^{-3}$ | $(2.9 \pm 1.4) \times 10^{-3}$ |
| $\eta_c \to 2(K^+K^-)$ | $(2.6^{+0.5}_{-0.4} \pm 0.4) \times 10^{-3}$ | $(1.5 \pm 0.7) \times 10^{-3}$ |
| $J/\psi \to \phi K^+K^-$ | $(1.2 \pm 0.3 \pm 0.1) \times 10^{-3}$ | $(1.83 \pm 0.24) \times 10^{-3}$ |
| $J/\psi \to 2(K^+K^-)$ | $(9.7^{+1.7}_{-1.6} \pm 1.0) \times 10^{-4}$ | $(7.8 \pm 1.4) \times 10^{-4}$ |

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