Enhancing the significance of gravitational wave bursts through signal classification

S Vinciguerra¹, M Drago², G A Prodi³, S Klimenko⁴, C Lazzaro⁵, V Necula⁴, F Salemi², V Tiwari⁶, M C Tringali³ and G Vedovato³

¹ University of Birmingham, Edgbaston, Birmingham, B15 2TT, United Kingdom
² Albert-Einstein-Institut, Max-Planck-Institut für Gravitationsphysik, D-30167 Hannover, Germany
³ Università di Trento, Dipartimento di Fisica, and INFN, TIFPA, I-38123 Povo, Trento, Italy
⁴ University of Florida, Gainesville, FL 32611, United States of America
⁵ INFN, Sezione di Padova, I-35131 Padova, Italy
⁶ Cardiff University, Cardiff CF24 3AA, United Kingdom

E-mail: serena.vinciguerra89@gmail.com

Received 28 November 2016, revised 10 February 2017
Accepted for publication 13 March 2017
Published 7 April 2017

Abstract

The quest to observe gravitational waves challenges our ability to discriminate signals from detector noise. This issue is especially relevant for transient gravitational waves searches with a robust eyes wide open approach, the so called all-sky burst searches. Here we show how signal classification methods inspired by broad astrophysical characteristics can be implemented in all-sky burst searches preserving their generality. In our case study, we apply a multivariate analyses based on artificial neural networks to classify waves emitted in compact binary coalescences. We enhance by orders of magnitude the significance of signals belonging to this broad astrophysical class against the noise background. Alternatively, at a given level of mis-classification of noise events, we can detect about 1/4 more of the total signal population. We also show that a more general strategy of signal classification can actually be performed, by testing the ability of artificial neural networks in discriminating different signal classes. The possible impact on future observations by the LIGO-Virgo network of detectors is discussed by analysing recoloured noise from previous LIGO-Virgo data with coherent WaveBurst, one of the flagship pipelines dedicated to all-sky searches for transient gravitational waves.
Keywords: gravitational waves, signal classification, machine learning, artificial neural networks, transient all-sky search

(Some figures may appear in colour only in the online journal)

1. Introduction: signal classification for background rejection

General searches for transient gravitational waves of generic waveform (GW *bursts*) have been accomplished exploiting the full sensitivity bandwidth of the laser interferometer gravitational-wave observatory (LIGO) [1, 2] and Virgo [3, 4] detectors. This type of *all-sky* search has been performed by analyzing the network detector data with coherent methods [5, 6] looking for signals lasting from *ms* to *s* scale. These methods successfully identified the first detected gravitational wave, GW150914 [7, 8]. In burst searches, the main factor which limits the statistical confidence of a gravitational wave candidate comes from non Gaussian noise outliers of single detectors, which may accidentally mimic a coherent response of the network. The implemented strategies to improve the capability of discriminating between signals and noise include both upstream and downstream methods. Upstream methods include data quality flags and vetoes at single detector level [9, 10] to clean the input of the network analysis. Downstream methods apply post processing procedures such as splitting the end results in a few separate frequency bands: to account for the most evident inhomogeneities of non Gaussian noise tails, any candidate belonging to a specific frequency band is thus ranked against the noise outliers characteristic of the same band.

Coherent WaveBurst (cWB) [11] is the flagship pipeline aiming at all-sky burst searches on LIGO-Virgo data using minimal signal assumptions. cWB has already been used for the analysis of data collected by the first generation of interferometers [12, 13] and during the first observation run of Advanced LIGO, O1. In September 2015 cWB was the first pipeline to identify GW150914 [8]. cWB is based on a likelihood maximization of the coherent response of the network, which also allows the reconstruction of the most significant signal characteristics [12, 14]. To further reduce the false alarm probability at a reasonable cost in terms of false dismissals, additional procedures have been implemented. These procedures include simpler tests, such as the rejection of candidate signals in case of unusually high energy disbalance at different detectors, as well as more elaborate methods. Among them, cWB uses procedures for constraining the polarization and direction of detectable signals, since noise spectra and directional sensitivities affect the fraction of detectors which significantly contribute to the coherent response of the network. All the strategies mentioned above preserve the degree of universality, typical of burst searches.

The performance of the pipeline can also take advantage of priors on the target signals; however, this is accomplished at the cost of a loss of generality of the search. Different versions of cWB pipeline have been tailored for, e.g. the search for coalescences of intermediate mass binary black holes [15] or of highly eccentric binary black holes [16].

The scope of this work is to demonstrate how signal classification methods can complement all-sky burst searches to enhance the significance of selected signal classes without losing the generality of the search. The proposed signal classification method is based on machine learning techniques [17] to identify such signals against the non Gaussian noise outliers recorded in LIGO-Virgo observations. The signal versus noise discrimination based on the use of machine learning techniques on single detector data has been discussed in several papers. For instance [18–20] address the classification of transient noise outliers and [21] the whitening of the detector output. A signal recognition approach based on boosted decision trees has been tested for the case of a burst search triggered by astrophysical events in a network of detectors [22].
Similarly, in [23], the authors developed a multivariate classification with random forests in the context of matched filtering searches for high-mass black hole binaries.

The novelty of our work lies in its integration into all-sky burst searches and on the implementation of new strategies for signal classification, taking advantage of pattern recognition in the framework of the time-frequency (TF) representation of the candidate signals. The latter is accomplished by averaging the response of more artificial neural networks (ANNs), as explained in section 2. As a case study, we focus on the recognition of signals consistent with coalescences of compact binaries (Section 3). Sections 4 and 5 discuss respectively the effectiveness of the ANN analyses of the TF representations and its robustness against changes in the target signal distribution. Section 6 explores the potentiality of a further multivariate step: we test a new ranking statistic built by combining the discrimination variables, adopted in a standard cWB analysis, with the output parameter describing the recognition of time-frequency patterns. By comparing the receiver operating characteristics, we show the achievable enhancement of GW burst searches through our signal classification strategies. In section 7, we give more general remarks about the impact of signal classification approaches on all-sky searches for GW bursts.

2. Methodology

With the purpose of enhancing the significance of GW bursts in all-sky searches, we complement the standard analysis pipeline cWB with a new discrimination variable based on ANNs, dedicated to signal classification according to their TF characteristics. In its standard operation, cWB uses two main post-processing statistics to discriminate gravitational waves from noise artifacts (glitches): the network correlation coefficient (cc) and the effective correlated SNR ($\rho$) [8]. The latter is used to rank candidate events and assign them a false alarm rate. The former measures the consistency of the candidate with a coherent response of the network to a GW. cWB uses TF transformations [24] at different resolutions (or levels) and defines the TF characterisation for each candidate, through e.g. a principal component analysis (PCA) [25]. We use a PCA since it summarises the main TF characteristics of each candidate in the least number of TF pixels from different resolutions.

To perform the ANN classification, the resulting TF representations in the whitened data domain are converted by a post-processing algorithm into a frame made by $8 \times 8$ pixels, each described by a scalar amplitude proportional to the fraction of total likelihood (see figure 1). Each frame carries no information about the absolute scale of the candidate strength, duration or frequency band and keeps only the uncalibrated shape information of the candidate’s TF trace (for more details see appendix A).

The 64 entries of the $8 \times 8$ frame feed the ANNs dedicated to the classification of the TF pattern. ANNs are informatic tools composed of calculation units (neurons, represented by circles in figure 2), connected together by synapses (represented by lines in figure 2), which acquire specific values (weights) accordingly to a supervised training procedure (see [26]).

---

7 In the algorithm dedicated to the definition of the frame, we also allow the rejection of late TF pixels with low likelihood. Starting from the pixels that appeared at the latest times, the procedure, which defines the $8 \times 8$ frame, discards a fraction of the TF pixels, corresponding to at most 10% of the total likelihood. This selection prevents the occurrence of problematic distortions of patterns in the $8 \times 8$ frames. Noisy pixels selected after the merger time would indeed shrink and move the characteristic trace left by chirping events on the $8 \times 8$ frame. This operation is particularly useful for the classification in real noise. In fact, it favours a neater TF map by rejecting any weak structures appearing after the merger time, which will then dominate the last column of the frame (see figure 1).

8 A supervised training procedure consists in an optimization rule tailored to minimize the error between actual and desired output over the selected target and complementary sets of signals.
The ANNs adopted for this study are multilayer perceptrons defined within ROOT [27] (the object oriented framework developed at CERN). Each ANN is composed of an input layer, fed by the $8 \times 8$ frame, and of 3 hidden layers, requiring the definition of about $2 \times 10^3$ synapses. Each ANN is trained against a set of $\sim 10^4$ target signals (type1 class) and a set of $\sim 10^4$ complementary events (type0 class) composed of either noise glitches or signals belonging to alternative classes. The classification rule consists in obtaining output values close to 1/0 for events classified as belonging to the target distribution, and close to 0 otherwise. Our procedure does not constrain the output value to be limited to [0,1] and overflows and underflows are possible.

**Figure 1.** Schematic representation of the proposed event classification. The procedure is based on the reconstructed time-frequency map of candidates. ANNs are trained to produce an output number close to 1 for events classified as belonging to the target distribution, and close to 0 otherwise. Our procedure does not constrain the output value to be limited to [0,1] and overflows and underflows are possible.

**Figure 2.** Schematic representation of the algorithm adopted to obtain the ANN average. In the square it is summarised the structure of each ANN: the black filled circles compose the input layer, while the empty ones represent elaborating units. Each of them sums the output values of the previous layer, by weighting them with the correspondent synapses’s (lines) strength. The new classification parameter is obtained by averaging the result of 4 independent ANNs (ANN average).

and references therein). The ANNs adopted for this study are multilayer perceptrons defined within ROOT [27] (the object oriented framework developed at CERN). Every ANN is composed of an input layer, fed by the $8 \times 8$ frame, and of 3 hidden layers, requiring the definition of about $2 \times 10^3$ synapses. Each ANN is trained against a set of $\sim 10^4$ target signals (type1 class) and a set of $\sim 10^4$ complementary events (type0 class) composed of either noise glitches or signals belonging to alternative classes. The classification rule consists in obtaining output values close to 1/0 for elements of type1/type0 class respectively. For more details on the structure and training procedure used for the present study, see appendix b and [27].
We mitigate the impact of the statistical fluctuations intrinsic to the ANN response by averaging the output of 4 (unless otherwise specified) independently trained ANNs and so introducing the ANN average, as shown in figure 2. Every considered ANN is built with the same fixed training procedure, but each of them is trained over an independent set of events sampled from the same distributions. The time required for the definition of a single ANN varies by several orders of magnitude according to the adopted training. However, once defined, ANNs are able to quickly evaluate cWB triggers, providing an effective discriminating variable, as illustrated in figure 3.

In this preliminary test, we considered a population of target signals (type1) made by compact binary coalescences (see table 2) and a population of accidental coherent responses of the initial LIGO-Virgo network of detectors (network glitches, type0). The former is produced by software injections of signals in a few days of initial LIGO and Virgo data from S6D-VSR3 runs [6, 9], recolored according to the early phase spectral sensitivity of advanced detectors [28]. The set of network glitches is produced by running the standard cWB all-sky search on the same recolored data streams after applying a set of time-shifts among different detectors so to cancel any physical correlation present. The use of actual data is necessary to model the non Gaussian noise features that dominate the performances of all-sky burst searches.

Figure 3 shows that ANNs can be effective in discriminating these target signals from network glitches, to a much more efficient degree than the more general criteria used by cWB.

3. Case study: classification of chirping signals

In this section we focus on the classification of a specific signal class, emitted by compact binary coalescences. The results will be highlighted in section 4 while their robustness will be discussed in section 5, proving that the enhancement provided by our classification still holds over a much wider signal parameter space than that used to train the ANNs.

3.1. Type 1 or target signal class

The results of the first observing run of advanced LIGO [29] and the measured rate for binary black hole (BBH) mergers of 9–240 Gpc$^{-3}$ yr$^{-1}$ [30] confirm that compact binary coalescences are the most numerous sources for current ground based detectors. In this study, we
focus on the classification of the inspiral phase of the binary evolution. During the inspiral, the gravitational emission is mainly determined by the chirp-mass $M_{\text{chirp}} = \frac{m_1 m_2}{m_1 + m_2}$, where $m_1$ and $m_2$ are the masses of the two companions. The chirp-mass drives:

- the frequency evolution $\dot{f} \propto M_{\text{chirp}}^{5/3} t^{1/3}$ of the gravitational waves, which sets the chirping behaviour commonly associated to these signals;
- the gravitational wave strain amplitude, which in time domain is: $\mathcal{A}(t) \propto \dot{f}^{2/3} M_{\text{chirp}}^{5/3}$;
- the time spent in the most sensitive part of the spectral sensitivity of the detectors, which scales as $t_{\text{det}} \propto M_{\text{chirp}}^{-5/3}$. This time is also strongly dependent on the noise spectral density (PSD) of detectors at lower frequencies.

Since the classification procedure depends on these properties, we investigated different distributions of chirp-masses. In particular, we mainly tested two signal populations:

- a **low-mass** distribution (see table 1) composed of signals characterised by a clear chirping feature in the TF representation, which is dominating the detectable signal within the spectral sensitivity of the detectors.
- a **wide-mass range** distribution (see table 2), composed of a CBC population including more massive systems and therefore shorter detectable signals, in which the inspiral phase plays a weaker contribution within the spectral sensitivity considered in this study (early phase of advanced detectors).

As CBC signals models, we adopt EOBNRv2 waveforms [32, 33]. They rely on the effective-one-body (EOB) formalism and describe all the phases (inspiral, merger and ring-down) of the coalescence. The actual distributions of the detected signals used in this study are a convolution between the signal population (from table 1 or 2) and the cWB detection efficiency. Figure 4 illustrates the selection effect due to this convolution on the wide-mass range signal distribution (table 2). As expected the detection pipeline is more sensitive to louder signals,

### Table 1. Main parameters of the low-mass CBC signal distribution.

| Parameter | Value |
|-----------|-------|
| Mass distribution | Uniform in $M_{\text{tot}}$ and $m_1 m_2$ |
| Total mass range ($M_{\odot}$) | $[3, 50]$ |
| Mass ratio range | $[1, 11]$ |
| Distance range $d$ (Mpc) | $\sim [70, 225]$ |
| Number of Shells$^a$ [31] | 3 |
| Distribution in each shell | Uniform in volume |

$^a$ The subdivision in shells is performed to decrease the computational load while ensuring more homogeneous statistical uncertainties on detection efficiency. The reference shell range is [100–150] Mpc and by rescaling signal amplitudes we populate the two contiguous shells ($\sim 70, 100$ Mpc and [150,225] Mpc).

### Table 2. Main parameters of the wide-mass range signal distribution.

| Parameter | Value |
|-----------|-------|
| Mass distribution | Uniform in $\log(m_1), \log(m_2)$ |
| Mass range ($M_{\odot}$) | $m_1, m_2 \in [1.5, 96.0], M_{\text{tot}} \in [3.0, 136.0]$ |
| Distance range $d$ (Mpc) | $\sim [45, 500]$ |
| Number of Shells$^a$ [31] | 6 |
| Distribution in each shell | Uniform in volume |

$^a$ As described in table 1, but rescaling signal amplitudes to populate five contiguous shells from $\sim 45$ Mpc to $\sim 500$ Mpc. Detection efficiencies are listed in table 3.
i.e. for more massive systems (figure 4). Indeed the gravitational wave amplitude scales as
$A(t) \propto t^{2/3} M_\text{chirp}^{2/5}$ and $M(M, q) = M_\text{chirp}^{3/5}(q + 1)^{-6/5}$, where $q = m_1/m_2$ is the mass ratio. Table 3 lists the overall detection efficiencies of cWB in the different shells. As expected, the pipeline efficiency decreases as the distance range increases.

### 3.2. Type0 or alternative signal class

Our case study requires to test the classification procedure with respect to both network glitches and alternative GW signals. The network glitch distribution has been described in section 2. As alternative signal class we considered a mixture of GW signal waveforms widely used in simulations of GW burst searches, i.e. the BRST set described in [6], which includes Gaussian pulses, sinusoidal signals with Gaussian amplitude envelope as well as White-Noise-Bursts waveforms. Such alternative signals lack of a proper astrophysical model and therefore their amplitude distribution at earth has been modeled by scaling their nominal amplitude value ($h_{\text{rss}} \equiv \sqrt{\int_{-\infty}^{\infty} h(t)^2 \, dt} = 2.5 \times 10^{-21}$) by a grid of logarithmically distributed scaling factors (0.075, 0.15, 0.3, 0.6, 1.2, 2.4, 4.8, 9.6, 19.2), see [6, 12]. All these signals have been injected on the same recolored data set used for the target signals injections.

### 4. Classification performance results

The main three tests of classification performances are summarized in Table 4.

The results are described in terms of the fraction of type1 (type0) events which are correctly classified (miss-classified) as belonging to the target class, $F_{i\rightarrow 1} (F_{i\rightarrow 0})$ defined by:
8

Here $N_k$ is the total number of tested events drawn from the $k$ class (where $k = 1, 0$ for type1, type0) and $\delta_i$ is defined for each event $i$ by:

$$\delta_i = \begin{cases} 1, & \text{if } \rho_i \geq \rho_{th} \land cc_i \geq cc_{th} \land \langle \text{ANN} \rangle_i \geq \langle \text{ANN} \rangle_{th} \\ 0, & \text{otherwise} \end{cases}$$

(2)

where the subscript ‘th’ refers to threshold value on the related variable. Events with $\delta_i = 1$ ($\delta_i = 0$) are classified as belonging to the type1 (type0) class. Figures 5–7 summarise the results of each test and provide a comparison of the effects of the standard test statistics of cWB (correlated SNR and network correlation coefficient) with the ANN average. Each figure is made by four plots, whose y-axis reports $F_{1\rightarrow1}$ and $F_{0\rightarrow1}$ respectively in the top ones and bottom ones. The plots to the left show $F_{1\rightarrow1}$ and $F_{0\rightarrow1}$ as a function of the threshold on the correlated SNR, $\rho_{th}$. The red lines represent results obtained without applying a threshold on the ANN average (or $\langle \text{ANN} \rangle_{th} = -\infty$), while green curves show the effects of selected threshold values, namely $\langle \text{ANN} \rangle_{th} = \{0.0, 0.5, 1.0\}$. The plots to the right show $F_{1\rightarrow1}$ and $F_{0\rightarrow1}$ as a function of the threshold on the ANN average. Here blue curves are computed by applying different thresholds on the correlated SNR $\rho_{th} = \{5.5, 6.7\}$. In all the plots a constant threshold on the network correlation coefficient is used, $cc_{th} = 0.6$, which is a common choice in standard GW burst searches [8]. This $cc$ threshold makes all the plotted fractions $F_{1\rightarrow1}$ and $F_{0\rightarrow1}$ lower than 1.

### 4.1. Test 1: low-mass chirp-like GWs versus glitches

The first classification test aims to discriminate network glitches from chirp-like signals of the low-mass distribution (first line of table 4). The results on a population of $5 \times 10^4$ events per each type are reported in figure 5: both correlated SNR and the ANN average are effective classifiers but their joint use shows advantages. In fact, it is possible to enhance the statistical confidence with a much smaller cost in terms of detection efficiency. For instance, considering the left plots, at any selected value of $\rho_{th}$, by adding $\langle \text{ANN} \rangle_{th} = 0.5$ almost the same fraction of chirp-like signals are recovered, while the fraction of mis-classified glitches $F_{0\rightarrow1}$ is reduced by about one order of magnitude. The right plots lead to similar considerations. $F_{1\rightarrow1}$ depends very weakly on $\langle \text{ANN} \rangle_{th}$ as long as $\langle \text{ANN} \rangle_{th} \leq 0.8$, while the mis-classified fraction of events $F_{0\rightarrow1}$ drops substantially in the same $\langle \text{ANN} \rangle_{th}$ range.

### 4.2. Test 2: low-mass chirp-like GWs versus alternative GWs

The signal classes considered for this test are the low-mass distribution and the BRST simulation set introduced as alternative signal class in 3.2. The results of the classification are
reported in figure 6, which shows the performances in terms of $F_{11}/\text{uni}$ and $F_{01}/\text{uni}$ with the same structure of figure 5. In particular, the right plots illustrate that the ANN average provides an efficient separation of the two GW populations, while the correlated SNR is instead agnostic with respect to the GW waveform class. In this test we selected a population of alternative GWs which are louder than the target GWs, to test an opposite condition with respect to what described in the previous subsection. The results are consistent with the ANN average being agnostic with respect to the loudness of the events, as it was designed to be. The chosen loudness of the alternative GW class carries no specific physical meaning, and the difference between $F_{11}/\text{uni}$ and $F_{01}/\text{uni}$ in the left plots cannot be used to classify type1 and type0 events in an actual GW search.

4.3. Test 3: wide-mass chirp-like GWs versus glitches

The task of the last classification test is separating network glitches from chirp-like signals, drawn from the wide-mass range simulation (table 2). We report the results in figure 7.

The plots of figures 5 and 7 exhibit very similar trends. The main difference is that $F_{11}$ values are generally slightly higher, since the wide-mass event distribution includes also louder GWs. Instead, even if the chirping character is weaker for the wide-mass signals distribution, the results show that the ANNs can be trained to give comparable performances with respect to the low-mass case.
5. Robustness

When the learning is supervised \[26\], as in our case, the ANN structures are defined thorough procedures mainly driven by the selected type0 and type1 samples. Since the astrophysical distribution of chirp-like GWs is unknown, we need to investigate the robustness of our approach against biases in the training distributions. To this purpose, we consider receiver operating curves (\(\text{ROC}\)) computed as \(F_{1\rightarrow1} / \text{versus} F_{0\rightarrow1} / \text{by varying the value of } \rho_{\text{th}}\), while keeping constant \(c_{\text{th}} = 0.6\), at selected ANN \(\text{values}.\) Figure 8 shows \(\text{ROC}\) curves for \(\langle \text{ANN} \rangle_{\text{th}} = -\infty\) and \(\langle \text{ANN} \rangle_{\text{th}} = 0.25\). The robustness is tested by comparing the improvement related to the application of the \(\langle \text{ANN} \rangle_{\text{th}} = 0.25\) on the \(\text{ROC}\) in two different cases:

- classification results are performed on chirp-like signals belonging to the same type1 distribution used for the training;
- classification results are performed on chirp-like signals belonging to a distribution different from that of the type1 used for the training.

Training set We build a simulation of GW emissions from three fixed equal-mass binaries \((10 \sim 10)M\sim (25 \sim 25)M\sim \) and \((50 \sim 50)M\sim \) (see table C1). From each of these three sources, we collect an equal number of reconstructed events to compose the training set for the target (type1) signals. The type0 training class is again defined with a sample of network glitches recoloured to mimic the spectral sensitivity in the early phase of the advanced detectors.
Testing sets

The first one is composed of samples of the same type1 and type0 populations used to define the training set just described, while the other one is formed by events drawn from the type1 and type0 classes of line 3 in table 4. For both cases, the tested type0 events are independent samples of the recoloured network glitches.

As expected, the ROC curves for a testing set different from the training set are worse, see figure 8. However, the improvement on the ROC achieved by implementing the threshold on the ANN average are significant and of comparable value in both testing sets. In this sense, figure 8 demonstrates the robustness of our approach against biases in the population-model. Since the mismatched type1 testing set has a much wider range of chirp-masses than the matched one, we present the achieved results of figure 8 as representative of the effect produced by increasing the assumed volume in the parameter space of the source population.

6. Multivariate analysis

The last versions of the cWB pipeline include an estimation of the binary chirp-mass $M_{\text{est}}$, computed as a best fit of the signal TF trace resulted from the PCA. To compute the fit, the algorithm considers all the pixels obtained with the TF decomposition at different levels, discarding the ones flagged as noise artifacts, according to the procedure defined in [8, 34]. Applying such post-processing analysis significantly improves the signal to noise discrimination achieved by cWB for CBC signals [34].

The joint distribution of the estimated chirp-mass and the ANN average (see figure 9) demonstrates that the two variables are not fully correlated. This gives the opportunity to further improve the results by a joint use of these variables. In addition we can consider other signal parameters estimated by cWB to implement a multivariate analysis (MVA).
For the morphological discrimination we are interested in, we implement an additional classification stage, using four MVA-ANNs (see figure 10). They are characterised by approximately 300 synapses and independently trained on $\sim 10^3$ events per class. The input quantities listed in figure 10 are then used to define a MVA-ANN average (\text{MVA}). We evaluated the effectiveness of the multivariate approach by comparing the ROC obtained from this MVA ranking statistic (MVA-ANN average) with the one driven by different values of correlated SNR, accordingly to the standard cWB analysis. Figure 11 clearly demonstrates that cWB’s performances in discriminating signals from glitches can be considerably improved by adopting the MVA-ANN average as ranking statistic. In fact, at fixed $F_{1\rightarrow 1}$ values $F_{0\rightarrow 1}$ is lowered down by three orders of magnitude, when switching the ranking statistic from the correlated SNR to the MVA-ANN average.

We finally performed checks to point out the sensitivity of our MVA-ANN average to the different inputs. These tests show that the estimated chirp-mass and the ANN average are by far the ones that impact more the results. Much smaller contributions comes from the central frequency, the correlated SNR and the network correlation coefficient, while duration and frequency bandwidth are the parameters which appear to be the least effective for our multivariate analysis.

7. Final remarks

In the previous sections we presented the application of the signal classification to the selected case study, i.e. CBC-like transient signals with an inspiral stage chirping up in frequency. The resulting enhancement of the significance of the detected signals in this
class is improved by orders of magnitude, at least in the confidence range investigated here (see figure 11). Alternatively, the gain can be interpreted as a significant increase of the detected fraction of sources at a given confidence, e.g. we recover $\sim 25\%$ more signals at a mis-classification of 0.01\% of noise events. The current approach could be further improved by considering more candidate parameters ($c_c$, $\rho$, $M_{\text{est}}$, etc), as proposed in section 6 by defining a new ranking statistic $MVA$ and by generalising the classification to other signal classes.

This search can be easily integrated within the framework of an all-sky search, where the search is general and open to every kind of GW-like signals with no particular assumptions on the morphology. In this situation, we can split the overall set of interesting triggers, $W$, into two classes, the say $A \subset W$ of CBC-like signals and the complementary $W-A$ of triggers belonging to $W$ but not to $A$. Following this approach for many waveform classes, all-sky searches can be managed as more separate searches on the same observation time, similarly to the more traditional case of all-sky searches performed on separate frequency bands. To account for the increased number of trials, a relative weight on these searches has to be chosen to portion out the overall false alarm probability of the all-sky search.

**Figure 9.** Distributions on plane defined by ANN average and the estimated chirp-mass of candidates belonging to the wide-mass signal class (top panel) and S6D recoloured glitches (bottom panel). The colour scale represents the fraction of recovered events, over a total of $5 \times 10^4$ per class. Negative values estimated chirp-mass are automatically set to zero.
Signal classes of astrophysical interest can be, for example, ring-down-like signals as emitted in quasi normal modes of NSs [35], say class B signals, or, more generically, signals with a reconstructed duration longer than a number of typical cycles, allowing for a diversity of waveform amplitude envelopes and phase evolutions in time, say class C signals. In addition,

Figure 10. Schematic representation of the algorithm adopted to define the MVA-ANN average. In the left square, we list all the input quantities elaborated by all the four ANN\textsubscript{MVA} to obtain \(\langle MVA \rangle\).

Figure 11. ROC using the MVA-ANN average (red) and the correlated SNR (blue) as a ranking statistic. In the last case, \(F_{1 \rightarrow 1}\) and \(F_{0 \rightarrow 1}\) are obtained considering a constant threshold on the network correlation coefficient \(cc_{th} > 0.6\) and ignoring the ANN average. The testing set, as well as the training one, has been defined by \(5 \times 10^4\) samples of S6D recoloured glitches (type0) and by \(5 \times 10^4\) signals drawn from the wide-mass range distribution (type1).

Signal classes of astrophysical interest can be, for example, ring-down-like signals as emitted in quasi normal modes of NSs [35], say class B signals, or, more generically, signals with a reconstructed duration longer than a number of typical cycles, allowing for a diversity of waveform amplitude envelopes and phase evolutions in time, say class C signals. In addition,
glitch classification methods can also be implemented for vetoing purposes, i.e. to reject the more frequent noise transient families at the detectors, say class Z. The most straightforward implementation of the all-sky search would then be a hierarchical signal classification scheme in subsequent steps, such as e.g. selecting \( A \subset (W - Z) \), then \( B \subset (W - Z - A) \), then \( C \subset (W - Z - A - B) \) and finally analysing the rest \( (W - Z - A - B - C) \). Such a hierarchical implementation would ensure that the considered classes are disjoint, both for signal and for glitches. Of course the more background rejection is accomplished in the first classification steps, the higher will be the resulting background left to the last ones.

In the framework of all-sky searches, the optimization strategy is not defined, both because we are lacking reliable source population models for most source classes and because we want to leave room for unexpected detections. The prioritization of signal classes and the portioning of the overall false alarm probability among the classes is subjective and has to be agreed upon, as a balance between boosting detection probability of better known sources and preserving suitable detection chances for signals in the widest accessible duration-frequency range. The former requires to take into account the detectable source number within the visible volume of the search; therefore, it would prioritize e.g. the frequency band of best spectral sensitivity and/or some waveform or polarization class. The latter instead calls for the consideration of the entire spectral range of the detectors, including disadvantaged spectral sensitivity bands, and for unmodelled waveforms and polarization states.

This issue has to be addressed anyway, regardless of the implementation of signal classification methods. In past all-sky searches, the portioning of the false alarm probability has been driven by uniform priors, i.e. by accounting a-posteriori for the trial factor coming from multiple sub-searches on different bandwidths or, in other words, by ranking signal candidates according to a quantity closely related to their inverse false alarm rate, as measured within the related sub-search. A portioning close to uniform makes sense also for the hierarchical search depicted here.

**Acknowledgments**

We would like to thank Ilya Mandel and Neil Cornish for the useful suggestions concerning this work.

The research leading to these results has received funding from the People Programme (Marie Curie Actions) of the European Union’s Seventh Framework Programme FP7/2007–2013/ (PEOPLE-2013-ITN) under REA grant agreement n ° [606176]. It reflects only the author’s view and that the Union is not liable for any use that may be made of the information contained therein.

**Appendix A. Conversion of TF representations into \( 8 \times 8 \) frames**

ANNs are here trained to recognise the common patterns in the TF representation of chirping signals. ANNs input layer are feed with the values of an \( 8 \times 8 \) frame constructed starting from the Time-Frequency Principal Component Analysis run by cWB 2G. The TF representation is defined by applying the Wilson–Daubechies–Meyer (WDM) transformation to the data [24] at different time (\( \Delta T \)) and frequency (\( \Delta F \)) resolutions (related to each other by \( \Delta T \times \Delta F = 1/2 \)). Then the algorithm applies the PCA to the most energetic pixels of each map (core pixels) to represent a particular event. From this multiple map representation, we analyse the possibility of discarding pixels, according to the rule mentioned in the first footnote 7. With this information we are able to focus (‘zoom’) on the TF region really involved by the event. This selected region is divided into fundamental units, i.e. the minimum time
resolution of the selected core pixels and half of the minimum frequency resolution (according to the application of WDM transformations). We finally group all the resulted units so to obtain a $\times 88$ frame. In order to adjust the ‘zoomed’ region to the $\times 88$ frame, each of the $\times 88$ squares can contain more fundamental units, or a fraction of them. The corresponding 64 values are therefore obtained by summing or spreading the likelihood of all fundamental units used to define each frame square. These values are then normalised and used to feed the ANNs. More details can be found in [27].

### Appendix B. Analysis and parameters

Table B1 reports the cWB-parameters adopted for the transient searches described in the paper. The training procedure, applied to construct the ANNs used for classifying TF patterns, is defined by the parameters reported in table B2 (for more details on the parameters and their choice see [37] and [27]).

### Appendix C. Other signal distribution

To test robustness of the proposed analysis against the uncertainty over the chirp-like signal distribution, we introduced another class of chirping signals, whose distribution is defined by the parameters reported in table C1.
References

[1] Aasi J et al 2015 Advanced Ligo Class. Quantum Grav. 32 074001
[2] Abbott B P et al 2009 LIGO: the laser interferometer gravitational-wave observatory Rep. Prog. Phys. 72 076901
[3] Acernese F et al 2014 Advanced Virgo: a second-generation interferometric gravitational wave detector Class. Quantum Grav. 32 024001
[4] Accadia T et al 2012 Virgo: a laser interferometer to detect gravitational waves J. Instrum. 7 P03012
[5] Abadie J et al 2010 All-sky search for gravitational-wave bursts in the first joint LIGO-GEO-Virgo run Phys. Rev. D 81 102001
[6] Abadie J et al 2012 All-sky search for gravitational-wave bursts in the second joint LIGO-Virgo run Phys. Rev. D 85 122007
[7] Abbott B P et al 2016 Observation of gravitational waves from a binary black hole merger Phys. Rev. Lett. 116 061102
[8] Abbott B P et al 2016 Observing gravitational-wave transient GW150914 with minimal assumptions Phys. Rev. D 94 069903
[9] Aasi J et al 2015 Characterization of the LIGO detectors during their sixth science run Class. Quantum Grav. 32 115012
[10] Abbott B P et al 2016 Characterization of transient noise in Advanced LIGO relevant to gravitational wave signal GW150914 Class. Quantum Grav. 33 134001
[11] Klimenko S et al 2016 Method for detection and reconstruction of gravitational wave transients with networks of advanced detectors Phys. Rev. D 93 042004
[12] Drago M 2010 Search for transient gravitational wave signals with unknown waveform in the LIGO-Virgo network of interferometric detectors using a fully coherent algorithm (http://paduaresearch.cab.unipd.it/2916/1/tesi_completa.pdf)
[13] Rahmanov M, Klimenko S, Mohanty S and Mitselmakher G 2005 Constraint likelihood analysis for a network of gravitational wave detectors Phys. Rev. D 72 122002
[14] Drago M 2009 Reconstruction of gravitational wave bursts with LIGO-Virgo network (https://dcc.ligo.org/public/0003/G0900548/005/Drago_Amaldi09_LIGO-G0900548-v5.pdf)
[15] Abadie J et al 2012 Search for gravitational waves from intermediate mass binary black holes Phys. Rev. D 85 102004
[16] Tiwari V et al 2016 Proposed search for the detection of gravitational waves from eccentric binary black holes Phys. Rev. D 93 043007
[17] Michalski R S, Bratto I and Kubat M 1999 Machine Learning and Data Mining: Methods and Applications (New York: Wiley)
[18] Troiano L, Rampone S, Pierro V and Pinto I M 2013 Neural network aided glitch-burst discrimination and glitch classification Int. J. Mod. Phys. C 24 1350084
[19] Powell J, Trifirò D, Cuoco E, Heng I S and Cavaglià M 2015 Classification methods for noise transients in advanced gravitational-wave detectors Class. Quantum Grav. 32 215012
[20] Zevin M et al 2016 Gravity spy: integrating advanced ligo detector characterization, machine learning, and citizen science (arXiv:1611.04596)
[21] De Rosa R, Eleuteri A, Milano L, Acernese F, Barone F and Tagliaferri R 2005 A hierarchical Bayesian framework for nonlinearities identification in gravitational wave detector outputs Class. Quantum Grav. 22 S1223
[22] Clark J, Adams T S, Meacher D, Jones G, Sutton P J and Minot A 2013 Gravitational-wave detection using multivariate analysis Phys. Rev. D 88 062006
[23] Hodge K A, Talukder D, Capano C, Baker P T, Caudill S and Cornish N J 2015 Multivariate classification with random forests for gravitational wave searches of black hole binary coalescence Phys. Rev. D 91 062004
[24] Klimenko S, Necula V and Mitselmakher G 2012 Transient analysis with fast wilson-daubechies time-frequency transform J. Phys.: Conf. Ser. 363 012032
[25] Vedovato G et al The cWB 2G clustering algorithms (https://atlas.aei.uni-hannover.de/~waveburst/doc/cwb/man/The-cWB-2G-clustering-algorithms.html#The-cWB-2G-clustering-algorithms)
[26] Stergiou C and Siganos D Neural networks (https://doc.ic.ac.uk/~nd/surprise_96/journal/vol4/cs11/report.html)
[27] Vinciguerra S 2014 cWB—artificial neural network (https://atlas.aei.uni-hannover.de/~waveburst/doc/cwb/man/ANN.html)

[28] Abbott B P et al 2016 Prospects for observing and localizing gravitational-wave transients with advanced LIGO and Advanced Virgo Living Rev. Relativ. 19 1

[29] Abbott B P et al 2016 Binary black hole mergers in the first Advanced LIGO observing run Phys. Rev. X 6 041015

[30] Abbott B P et al 2016 The rate of binary black hole mergers inferred from Advanced LIGO observations surrounding GW150914 Astrophys. J. Lett. 833 L1

[31] Mazzolo G 2013 Search for intermediate mass black hole binaries with networks of ground-based gravitational-wave detectors PhD Thesis Observational Relativity and Cosmology, AEI-Hannover, MPI for Gravitational Physics, Max Planck Society

[32] Boyle M, Buchman L T, Kidder L E, Pfeiffer H P, Pan Y, Buonanno A and Scheel M A 2011 Inspiral-merger-ringdown multipolar waveforms of nonspinning black-hole binaries using the effective-one-body formalism Phys. Rev. D 84 124052

[33] Shibata M, Tagoshi H, Mino Y, Sasaki M and Tanaka T 1997 Black hole perturbation Prog. Theor. Phys. Suppl. 128 1–121

[34] Necula V, Tiwari V, Klimenko S and Mitselmakher G 2015 Reconstruction of chirp mass in searches for gravitational wave transients Class. Quantum Grav. 33 01LT01

[35] Gusakov M E, Gualtieri L, Kantor E M and Chagunov A I 2014 Quasinormal modes of superfluid neutron stars Phys. Rev. D 90 024010

[36] Mercer A, Klimenko S, Yakushin I and Mitselmakher G 2008 A coherent method for detection of gravitational wave bursts Class. Quantum Grav. 25 114029

[37] ROOT 2016 Tmultilayerperceptron class reference (https://root.cern.ch/doc/master/classTMultilayerPerceptron.html)