VARIATIONS IN HIGH DEGREE ACOUSTIC MODE FREQUENCIES OF THE SUN DURING SOLAR CYCLES 23 AND 24

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ABSTRACT

We examine continuous measurements of the high-degree acoustic mode frequencies of the Sun covering the period from 2001 July to 2014 June. These are obtained through the ring-diagram technique applied to the full-disk Doppler observations made by the Global Oscillation Network Group. The frequency shifts in the degree range of 180 ≤ ℓ ≤ 1200 are correlated with different proxies of solar activity, e.g., 10.7 cm radio flux, the International Sunspot Number, and the strength of the local magnetic field. In general, a good agreement is found between the shifts and activity indices, and the correlation coefficients are found to be comparable with intermediate-degree mode frequencies. Analyzing the frequency shifts separately for the two cycles, we find that cycle 24 is weaker than cycle 23. Since the magnetic activity is known to be different in the two hemispheres, for the first time, we compute the frequency shifts over the two hemispheres separately and find that the shifts also display hemispheric asymmetry; the amplitude of shifts in the northern hemisphere peaked during late 2011, more than two years earlier than in the south. We further correlate the hemispheric frequency shifts with the hemispheric sunspot number and mean magnetic activity index (MAI). Since the frequency shifts and the hemispheric activity indices are found to be significantly correlated, we suggest that the shifts be used as an indicator of hemispheric activity since not many indices are measured over the two hemispheres separately. We also investigate the variation at different latitudinal bands and conclude that the shifts in active latitudes correlate well with the local MAI.

Key words: Sun: activity – Sun: atmosphere – Sun: helioseismology – Sun: magnetic fields – Sun: oscillations

1. INTRODUCTION

The eigenfrequencies of solar oscillations that provide insight into the interior structure and dynamics of the Sun are known to vary with solar activity (Woodard & Noyes 1985; Elsworth et al. 1990; Woodard et al. 1991). With improved observations, the relationship between mode frequencies and activity cycles has been extensively studied and it is well-established that the frequencies exhibit a good correlation with many different solar activity proxies (Chaplin et al. 2007; Jain et al. 2009). Comprehensive analyses using consistent and longer data sets both from ground- and space-based instruments, however, have shown different and unexpected results. For example, a phase-wise analysis of the intermediate-degree frequencies showed that the degree of correlation is different from phase to phase of the cycle (Jain et al. 2009). The extended minimum phase between cycle 23 and cycle 24 also revealed significant differences between acoustic frequencies and solar activity (Broonhall et al. 2009; Salabert et al. 2009; Tripathy et al. 2010b). Analyzing Birmingham Solar Oscillations Network data over three solar cycles, Basu et al. (2012) presented evidence that the solar sub-surface layer was very different between cycles 22 and 23, with deeper changes occurring during cycle 22 compared to cycle 23. A recent analysis of low-degree modes from the Global Oscillations at Low Frequency instrument also showed that cycle 24 is magnetically weaker in the upper sub-surface layers of the Sun (Salabert et al. 2015).

Since the frequency shifts are predominantly a function of the frequency alone, it has been suggested that the perturbations causing the changes are confined to the shallow layers (Elsworth et al. 1994; Howe et al. 1999; Chaplin et al. 2001; Simoniello et al. 2013b). In addition, a small but statistically significant change in modes with turning points at or below the convection zone has also been reported (Chou & Šerebryanskij 2005; Baldner & Basu 2008). Thus, the physical mechanism still remains an open question, i.e., whether the shifts are caused by structural changes in the sub-surface layers, such as changes in temperature (Kuhn 1988), or by the magnetic fields either through the action of the Lorentz force on the plasma or indirectly by modifying the physical properties, e.g., the size of the acoustic cavity (Dziembowski & Goode 2005). Another possibility is the influence of the magnetic atmosphere (Jain & Roberts 1996). The analysis of frequency shifts by Dziembowski & Goode further suggest that it is the indirect effect that dominates, particularly changes to the near surface stratification resulting from the suppression of convection by the magnetic field. On the other hand, they also reported small but significant departures from the simple scaling of the lower-frequency p-mode frequency shifts with the inverse mode inertia, implying that there is a contribution from deeper layers due to the direct effect of the magnetic fields. However, through theoretical modeling Foulon & Roberts (2005) concluded that the effect of a buried magnetic field on the frequencies of solar acoustic modes is too small to account for the changes in frequency over the solar cycle. Moreover, Howe et al. (2002) showed that the temporal and latitudinal distribution of the frequency shifts is correlated with the spatial distribution of the surface magnetic field, implying that the shifts cannot be purely explained by structural changes. Thus, a thorough modeling is required to better understand the cause of the frequency shifts.

In this context, it is desirable to analyze high-degree modes whose lower turning points occur very close to the solar surface (Figure 1). But the determination of high-degree mode frequencies is a challenging task since individual modes blend into ridges at high degrees (see, for example, Korzennik et al. 2013b, and the references therein). In spite of this inherent
difficulty, there have been several attempts to compute and analyze high-degree oscillation frequencies. Using time series from the Michelson Doppler Imager (MDI) on board the Solar and Heliospheric Observatory (SOHO), Rhodes et al. (2002) estimated the high-degree mode frequencies and demonstrated the response of the modes to short-term changes in solar activity. Rose et al. (2003) extended this analysis to include data from the Global Oscillation Network Group (GONG). Considering global modes from MDI, Rabello-Soares et al. (2008b) concluded that the changes in the high-degree mode frequencies are in good agreement with the intermediate-degree results, except for years when the MDI was highly defocused. A similar study, but including more data, demonstrated a quadratic relation between the frequency shifts and activity indices (Rabello-Soares 2011), which was not confirmed by Tripathy et al. (2013b). Rhodes et al. (2011) also investigated the temporal changes in high-degree p-mode frequencies and pointed out that short-duration, high-frequency shifts are more sensitive to changes in solar activity than the shifts measured by long-duration studies consisting only of the lower frequencies. Recently, Korzennik et al. (2013a) computed high-degree modes for three different epochs from three different instruments and compared the mode parameters.

Ring-diagram analysis (Hill 1988), in addition to global helioseismic techniques, measures the characteristics of high-degree modes (180 ≤ ℓ ≤ 1200) in localized regions and has been primarily used to investigate the variations around active regions (Rajaguru et al. 2001; Bogart et al. 2002; Rabello-Soares et al. 2008a). Hindman et al. (2000) found that the frequencies act as a tracer of local magnetic activity, which was later shown to be less reliable during low-activity periods (Tripathy et al. 2011). However, the response of the frequencies of the high-degree modes to solar activity using a continuous and consistent data set has not been undertaken previously. The results reported earlier mostly use discrete data sets obtained by the MDI instrument when it was operating in its high-resolution mode for three months every year. Our goal in this paper is to investigate the behavior of high-degree mode frequencies over a period of about 13 years consisting of both the descending phase of cycle 23 and the ascending phase of cycle 24. This analysis supplements our previous study (Tripathy et al. 2013a, hereafter paper I) where we primarily focused on the minimum phase of solar cycle 23. Since magnetic activity (e.g., the sunspot number and polar flux) exhibit north–south hemispheric asymmetry, for the first time we examine the variations in oscillation frequencies over these two hemispheres separately. We further explore the relationship between the frequency shifts and magnetic activity at different latitudes in both hemispheres.

The rest of the paper is organized as follows. The data and the analysis procedure are described in Section 2, while the results are described in Section 3. Finally, we summarize the main conclusions in Section 4.

2. DATA ANALYSIS

The GONG has been operating in high-resolution mode since 2001 July and produces full-disk Doppler images of the Sun every minute. These images are then processed through the ring-diagram pipeline (Corbard et al. 2003) to produce the mode frequency tables. In short, the ring-diagram technique (Hill 1988) produces a three-dimensional spectral (two in space, one in time) of a small region on the surface of the Sun tracked and remapped over a given time period. In the GONG pipeline, each region (hereafter tile) covers an area of 15° × 15° in heliographic longitude and latitude and is tracked for a period of 1664 minutes (hereafter ring-day). The centers of the tiles are separated by 7°5 and a set of 189 overlapping tiles, referred to as dense-pack, covers the region of ±52°5 in latitude and central meridian distance from the disk center. The three-dimensional spectra is fitted with a Lorentz profile model (Haber et al. 2000)

\[
P(k_x, k_y, \nu) = \frac{A}{(\nu - \nu_0 + k_x U_x + k_y U_y)^2 + \Gamma^2} + b/k^3,
\]

where \(P\) is the oscillation power for a wave with a temporal frequency \(\nu\) and the total wave number \(k^2 = k_x^2 + k_y^2\). There are six parameters to be fitted: the zonal and meridional velocities, \(U_x\) and \(U_y\), respectively, the central frequency, \(\nu\), the amplitude, \(A\), the mode width, \(\Gamma\) and the background power, \(b\). For each tile, the peak fitting procedure returns about 200 modes with degree, \(\ell\), between 180 and 1200 and with frequencies between 1700 and 5600 \(\mu\)Hz. However, it is to be noted that these are not discrete modes in the same sense as those observed at lower degrees. The ring frequencies are taken from the GONG website\(^1\) and cover the period from 2001 July 26 to 2014 June 19.

Since GONG is a ground-based network of six stations, poor observing conditions and instrument failures give rise to gaps in the data. These gaps are measured in terms of a fill factor, where a fill of unity means no data loss (100% coverage). The distribution of the fill factor of the ring-data is shown in Figure 2. Since gaps in the data cause biases and data of lower fill may not be fitted reliably, we exclude data where the fill factor is lower than 70%. This resulted in 3761 out of a total of 4146 ring-days and consists of approximately 710,800 tiles.

The mode frequencies (and other mode parameters) measured by the ring-diagram technique are affected by

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\(^1\) http://nisp.nso.edu
foreshortening as well as gaps in observation. Following the method outlined by Howe et al. (2004), we model the effects of position of the disk as a two-dimensional function of the distance from disk center ρ without the cross-terms, combined with a linear dependence on the fill factor, f(t),
\[
\nu(\rho_x, \rho_y, t) = a_0 + a_1 \rho_x + a_2 \rho_y^2 + a_3 \rho_y + a_4 \rho_x^2 + a_5 f(t),
\]
where \(\rho_x\) and \(\rho_y\) are the longitudinal and latitudinal component of \(\rho\) and \(a_i\) are coefficients determined by fitting. In order to minimize the effect of the magnetic field on the fitting, we use the frequencies over a two-year period when solar activity is low (between 2007 September 1 and 2009 August 31) and evaluated the \(a_i\)s for each \((n, \ell)\) mode. The resultant coefficients and the resultant differences between the corrected and original frequencies are shown in Figures 3 and 4 of paper I.

Figure 3 compares the fitted modes obtained from global and ring-diagram analysis techniques and data sources. The black circles display the fitted intermediate-degree modes using MDI data, the red dots represent the global high-degree modes from GONG, and the blue circles display the high-degree modes obtained through the ring-diagram analysis of the GONG data.

Figure 4. Frequency differences between 2002 and 2010. The plus symbols denote the differences computed from high-degree global modes obtained from Korzennik’s fitting in the degree range of 100 and 1000. The filled circles denote the differences between high-degree modes computed using the ring-diagram technique. The 2002 frequencies are calculated from a time series of 98 days, with a start day of 2002 February 14, while those in 2010 are computed from a time series of 67 days, starting from 2010 May 7.

Figure 5. Temporal evolution of the mean magnetic activity index, MMAI. The solid black line shows the running mean over 181 points. The dashed red line represents the 181 day smoothed 10.7 cm radio flux measurements scaled linearly with MMAI.

2 https://www.cfa.harvard.edu/bsylvain/research/tables/HIL
appear to fall along the same ridges. This also suggests that the systematic errors in the frequencies arising from analysis choices are independent of time.

We further compare the individual frequencies obtained from GONG data between high — (2002) and low — (2010) activity periods where the frequencies corresponding to low-activity period are computed from a 67 day time series starting from 2010 May 7 (Korzennik et al. 2013a). Figure 4 shows the frequency differences. The pluses denote the differences between the global high-degree modes while the filled circles show the differences between the modes obtained from the ring-diagram analysis. The mode frequencies obtained from two different methods agree very well except for the scatter seen in the global modes at high frequencies.

For the purpose of characterizing activity associated with each tile, we compute a magnetic activity index (MAI) that indicates the local magnetic field strength associated with each tile. The index is calculated by tracking and remapping the line-of-sight magnetograms during the same interval as the Dopplergrams and averaging the absolute values of all of the pixels with a field strength higher than a threshold value. Since the entire period under study is not covered by the MDI

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**Figure 6.** (a): Pearson’s linear correlation coefficient, $r_p$, between $\delta
\nu_i$ and MAI over 189 dense-pack tiles for each ring-day as a function of the mean MAI. The symbols and colors represent the year-wise variation and are indicated in the top panel. Panels (b)–(e) show the coefficients around the maximum (2001–2003; 2012–2014), descending (2004–2006), minimum (2007–2009), and ascending (2010–2011) phases of the solar activity in cycles 23 and 24, respectively. Note that the range of the mean MAI in each panel is different.
magnetograms alone, we use the 96 minute MDI magnetograms up to 2007 April 9 and GONG magnetograms sampled every 32 minutes thereafter. Lower threshold values of 50 and 8.8 G are used for MDI (Basu et al. 2004) and GONG, respectively, which are approximately 2.5 times the estimated noise in the individual magnetograms. To generate a uniform MAI, the MDI MAI values are scaled to GONG values using a conversion factor of 0.31, which was obtained by relating the line-of-sight magnetic field measurements between these two instruments using a histogram-equalizing technique (Riley et al. 2014).

A well-known systematic effect in solar physics, especially with ground-based resolved data, is the annual variation with the solar inclination toward Earth, which can lead to a loss in spatial resolution and hence change the geometric foreshortening. This effect is clearly seen in the MAI, the frequency shifts (calculated later), as well as in solar activity indices. We have applied a technique based on Fourier Transform to remove the 1 year periodicity from all the date sets used in the study of temporal variations. However, for a better demonstration, uncorrected data has been included in some figures by means of symbols. Also note that the frequency shifts and MAIs have been interpolated to 4146 ring-days before the Fourier Transform was applied.

Figure 5 shows the temporal variation of the mean MAIs (MMAIs) where the mean is taken over 189 dense-pack tiles of each ring-day. We also calculate a 181 point running mean, which approximately corresponds to seven Carrington rotations in terms of ring-days, to smooth out the short-term variations. This is shown by solid (or dashed) lines in most of the figures. In the same plot, we also show the variation of the 10.7 cm radio flux, $F_{10.7}$, as a measure of the solar activity (dashed red line). Significant agreement is seen between the shape of the curves, indicating a high correlation, which is confirmed with a rank correlation coefficient value of 0.97. Thus the MMAIs will be considered as an activity proxy in the rest of the analysis.

In order to examine the sensitivity of the frequency shifts to solar activity, we further use two commonly used activity proxies; the International sunspot Number, ISN, and $F_{10.7}$, both obtained from the National Geophysical Data Center (NGDC). The sunspot number is a weighted estimate of the individual sunspots and sunspot groups derived from inspection of the solar photosphere in white light, while the radio flux index is a measurement of the integrated emission at the 10.7 cm wavelength from all sources present on the solar disk.

Figure 7. Histogram showing the distribution of the correlation coefficients as a function of the number of days.

3. RESULTS AND DISCUSSIONS

3.1. Spatial Variation

Hindman et al. (2000) studied the spatial variation of the frequencies obtained from ring-diagram analysis and showed that the active regions appear as locations of large frequency shifts. This gave rise to the conclusion that the frequency shifts can be used as a tracer of surface magnetic activity. Subsequent analysis (Tripathy et al. 2010a; Tripathy et al. 2011), however, indicated that the conclusion is valid only during the period of high activity and breaks down during the activity minimum period.

We estimate the spatial frequency shift, $\delta v_s$, weighted by the fitting error, $\sigma_{n,t}$, for each tile according to

$$\delta v_s(t) = \frac{1}{\sum_t \sigma_{n,t}^2} \left( \sum_t \frac{\delta v_{n,t}(t)}{\sigma_{n,t}^2} \right),$$

where the frequency difference of each mode, $\delta v_{n,t}$, is calculated with respect to the error-weighted average of that particular mode present in the dense-pack of each ring-day. Sequences of images showing $\delta v_s$ and MAI for four days during the high- and low-activity periods are presented in Figures 1 and 2 of Tripathy et al. (2011), demonstrating the relationship between them.

http://www.ngdc.noaa.gov/stp/solar/solardataservices.html
For a quantitative analysis, we calculate the Pearson’s linear correlation coefficient, $r_p$, between $\delta\nu$ and MAI of each tile in the dense-pack. The resultant coefficients (color coded for different years) as a function of the MMAIs are shown in Figure 6(a). We observe that for high values of MMAI the correlation is stronger, while for MMAI $\leq$ 1 G, the correlation coefficients are small, with a median value of 0.41. Upon closer inspection, we also find higher correlation coefficients for small MMAI’s, especially during the minimum phase. For a better visualization, we plot the correlation coefficients for different phases separately. Figure 6(b) shows the period around the maximum phase of cycles 23 (2001–2003) and 24 (2012–2014) and illustrates a stronger correlation between the frequency shifts and mean MAI. The coefficients for the descending phase of cycle 23 (2004–2006), the extended minimum phase (2007–2009), and the ascending phase of cycle 24 (2010–2011) are shown in panels (c), (d), and (e) of Figure 6, respectively. In all cases, we
note both high and low correlation coefficients for MMAI values of \(\approx 1\) G and a large spread in the values of the coefficients (largest for the extended minimum phase). We further observe negative coefficients in panels (c) and (d) (fewer in panel (d)), implying anti-correlation, but the values are not significant. Analysis of global mode frequencies with solar activity proxies during the extended minimum phase between cycles 23 and 24 has also indicated anti-correlation between them (Salabert et al. 2009; Tripathy et al. 2010b). Thus our result is consistent with the global low- and intermediate-degree modes.

The distribution of the linear correlation coefficient in the form of a histogram is shown in Figure 7. Only a small fraction of 63 ring-days (1.7%) are anti-correlated with MMAI, while 73% and 65% ring-days have positive correlation coefficients that are higher than 70% and 80%, respectively.

Figure 9. Correlations between the frequency shifts and different activity indicators for the cycle 23 descending (filled circles) and cycle 24 ascending (open circles) phases. The straight lines are the linear regressions for cycle 23 (blue), cycle 24 (green), and all data (red). The activity proxies are given on the abscissae. The black line is the result of a quadratic fit using all data.
Table 2
Gradients and Linear Correlation Coefficients of Frequency Shifts as a Function of Activity Proxies

| Activity Index | Cycle 23 Gradient $\pm$ $10^{-4}$ | $r_p$ | Cycle 24 Gradient $\pm$ $10^{-4}$ | $r_p$ | All Data Gradient $\pm$ $10^{-3}$ | $r_p$ |
|---------------|---------------------------------|------|---------------------------------|------|---------------------------------|------|
| MMAI          | $0.355 \pm 1.58 \times 10^{-4}$  | 0.91 | $0.242 \pm 1.60 \times 10^{-4}$  | 0.90 | $0.289 \pm 1.09 \times 10^{-1}$ | 0.88 |
| $F_{10.7}$    | $0.009 \pm 3.99 \times 10^{-6}$  | 0.87 | $0.005 \pm 3.14 \times 10^{-6}$  | 0.79 | $0.009 \pm 3.45 \times 10^{-3}$ | 0.87 |
| ISN           | $0.009 \pm 4.22 \times 10^{-6}$  | 0.87 | $0.009 \pm 4.89 \times 10^{-6}$  | 0.79 | $0.009 \pm 3.45 \times 10^{-3}$ | 0.87 |

**Note.** Units of the gradients are $\mu$Hz per unit activity, where the activity units are defined in Table 1.

Figure 10. Temporal variation of (a) $\delta \nu$, (b) magnetic activity, and (c) hemispheric sunspot number (HSN) for the northern hemisphere (blue circles) and the southern hemisphere (orange circles). The solid line in each panel represents the running mean over 181 points, the blue and red lines denoting the northern and southern hemispheres, respectively. The dashed and dotted–dashed lines indicate the peak amplitudes in the northern and southern hemispheres, respectively.
3.2. Variation with Solar Cycle

The frequency variation with the solar cycle is investigated by calculating a mean frequency shift, $\delta \nu$, using a formula that is generally applied to compute frequency shifts of global modes,

$$\delta \nu = \sum_{n, \ell} \frac{1}{\sigma^{2}_{n, \ell}} \sum_{i=1}^{189} \frac{\delta \nu_{n, \ell}(i)}{\sigma^{2}_{n, \ell}},$$

where $\delta \nu_{n, \ell}$ is the frequency difference of each $(n, \ell)$ mode with respect to the average frequency of the same mode over the dense-pack tiles corresponding to a quiet ring-day 2011 May 2008. Table 1 shows the values attained by the mean frequency shifts and three activity proxies during cycles 23 (up to 2008 December 31) and 24. However, the values presented here may not be the true minimum and maximum values of the solar cycles (and hence the ratio) due to the period considered in this study. The final column in the table shows the ratio of amplitudes of cycle 23 to cycle 24, in the sense $(\text{max} - \text{min})_{23} / (\text{max} - \text{min})_{24}$. While the mean shifts and activity indices show a higher amplitude in cycle 23, the MMAI, which is the strong field component of the magnetic flux, shows a higher amplitude in cycle 24. This is also evident in Figure 8, where $\delta \nu$ and activity proxies are plotted as a function of time. The solid line in each panel represents the running mean over 181 points and guides the eye to visualize the descending phase of cycle 23, ascending phase of cycle 24 and the extended minimum period between the two cycles. While the extended minimum phase is unambiguously detectable in the activity indices, it is not noticeable in the frequency shifts illustrating that (i) the change in frequency with activity is no longer linear during the solar minimum and (ii) the minimum phase in frequency lasted over a shorter period compared to the solar activity. This is consistent with our earlier conjecture that the changes in high-degree frequencies may be associated with two components: strongly localized active regions and the small-scale horizontal magnetic field (Tripathy et al. 2010a). The former explains the larger shifts, while the latter is responsible for the smaller shifts in all phases of the solar activity, including the minimum phase. We return to this discussion in Section 3.4.

We further examine the sensitivity of the frequency shifts to solar activity by performing a linear fit between them using data over the entire 13-year period as well over different phases separately (Figure 9). The straight lines are the result of linear regressions for cycle 23 (blue), cycle 24 (green), and all data (red). The black solid line in each of the panels is the result of a quadratic fit using all data and differs from the linear fit (red line) when activity is high. Thus, we conclude that the activity proxies display simple linear relations with the frequency shifts instead of a quadratic relation as reported earlier (Rabello-Soares 2011). We conjecture that the quadratic relation is an outcome of using fewer data points as we observe at a high-activity period. Table 2 shows the frequency shift per unit activity gradients and the corresponding linear correlation coefficients for each of the proxies. For MMAI, the gradient is higher for cycle 23 compared to cycle 24, while for other proxies the gradients do not vary significantly between different cycles or over the entire time span. The level of correlation with different activity proxies is higher than 80%, largely in good agreement with the analysis of intermediate-degree global modes (Jain et al. 2011); there are, however, marginal changes between cycles 23 and 24, and the entire period. Also, as anticipated, the correlation is higher between $\delta \nu$ and magnetic activity measured locally.

Before the data were corrected for the annual variation, a secondary periodicity, except during the minimum phase, was observed in the frequency shifts as well as in the activity indices. This signal is attributed to the quasi-biennial signal (Broomhall et al. 2012; Simioniello et al. 2013b) that has recently received a great deal of interest but is beyond the scope of this paper.

3.3. Hemispheric and Latitudinal Variation

In addition to the temporal changes, magnetic activity is generally north–south asymmetric and varies with latitudes. The hemispheric asymmetry in sunspot number and polar flux is well-established and provides important clues for understanding the mechanism of solar cycle (for a review see, Norton et al. 2014). The variation can be easily seen in magnetic butterfly diagrams where the activity belts move toward the geometric equator with the progression of the solar cycle. In addition, the wings of the butterfly diagram show a high degree of mirroring across the equator, indicating coupling between the two hemispheres. In this context, the frequencies obtained from the ring-diagram technique provide a natural way to investigate variations between different hemispheres as well as latitudes, since the $m$-averaged global frequencies of oscillation modes, commonly used for solar cycle studies, are latitudinally averaged and cannot distinguish between the northern and southern hemispheres. Nonetheless, the frequency of individual

| Activity Index | Min | Max  | Gradient                      | $r_p$ | $r_s$ |
|----------------|-----|------|------------------------------|-------|-------|
| MMAI (G)      | 0.04| 8.09 | Northern Hemisphere          | 0.285±1.61×10^{-4} | 0.90 (0.97)² | 0.90 (0.97)² |
| HSN           | 0   | 110  | Southern Hemisphere          | 0.022±1.32×10^{-5} | 0.84 (0.98) | 0.84 (0.97) |
| MMAI (G)      | 0.00| 0.00 |                             | 0.284±1.34×10^{-3} | 0.86 (0.94) | 0.87 (0.95) |
| HSN           | 0   | 126  |                             | 0.022±1.09×10^{-5} | 0.86 (0.98) | 0.85 (0.98) |

Note:
² The calculated probabilities of having null correlations are 0.
³ The values in the brackets are the correlation coefficients corresponding to the 181 point running mean.
multiplets (which are function of $n$, $\ell$, $m$) can be used to study latitudinal variations based on $m/\ell$ values (e.g., Howe et al. 2002; Jain et al. 2011). The sectoral and near-sectoral modes are sensitive to the region near the equator while the zonal and tesseral modes are sensitive to the wider range of latitudes.

Figure 10 shows the mean frequency shifts for northern (blue circles) and southern (orange circles) hemispheres and the solid line represents the running mean over 181 points. The shifts are calculated using Equation (3), where $\delta n_{n,\ell}$ corresponds to the frequency difference of each mode with respect to an average frequency over the same hemisphere. Both the northern and southern frequency shifts decreased starting from mid-2001 and their variations are similar until the end of 2007. Around the beginning of cycle 24, the two shifts diverged and reached their maximum amplitude at two different periods. In cycle 24,
the northern shifts peaked in late 2011, more than two years earlier than the south, as indicated by the dashed and dotted lines, respectively. This closely agrees with the hemispheric asymmetry seen in other activity indices, e.g., the MAI (panel (b)) and the hemispheric sunspot number (HSN, panel (c)), where the latter is obtained from the NGDC. Conspicuously, both the dashed and dotted-dashed lines are aligned between different panels, illustrating that the frequency shifts very closely follow the changes in activity indices over both hemispheres. To quantify the relationship, we compute the correlation coefficient between them and the shift per unit activity gradients. These are tabulated in Table 3. The relative change in the gradients points out that the shifts remained approximately constant between the two hemispheres. Moreover, the Pearson’s linear and Spearman’s rank correlation coefficients further confirm the linear relation between the shifts and activity indices. The values mentioned in the brackets are the coefficients corresponding to the smoothed values, and as expected, signify a closer interaction. In the absence of many activity proxies for different hemispheres, we suggest that the frequency shifts may be used as a hemispheric activity proxy.

Figure 11 shows the mean frequency shifts for a few selected latitudes where the shifts are obtained by subtracting the average frequency computed at the same latitude covering all data sets but only over the central meridian. For clarity, we only show the mean shifts resulting from a running mean over 181 points. The shifts show significant variations between different latitudes as well as between the two cycles. Moreover, a clear extended minimum phase is not seen at any of the latitudes. In cycle 23, peak amplitudes are observed around the beginning of 2002 for latitudes $7.5^\circ$ N and $15^\circ$ N, while for southern hemisphere the peak amplitude can be seen for all the latitudes around the same time. On the contrary, we notice opposite behavior in cycle 24; the peak amplitude is seen around the end of 2011 for all latitudes, while in the southern hemisphere the shifts peak around the beginning of 2014, except for at $30^\circ$ latitude. At this latitude, a peak amplitude is observed around 2012 November and a decreasing trend afterward. Interestingly, the shifts corresponding to the equatorial belt have only one strong peak amplitude in 2002 and a highly reduced amplitude (about half of cycle 23) for cycle 24. Thus, if we consider frequency shifts that correspond to the equator we will conclude that cycle 24 is weaker compared to cycle 23, while at other latitudes, the shifts are nearly similar (slightly higher for southern hemisphere) between cycles 23 and 24.

The frequency shift per unit of magnetic activity (filled circles) and the Pearson’s linear correlation coefficients (filled squares) between MAI and $\delta v$ for different latitudes are shown in Figure 12. Higher values of the gradients are found, corresponding to the solar active latitude belt, i.e., about $\pm 30^\circ$. In the same latitudes, magnetic activity as measured by MAI are also found to be larger and hence a strong correlation is observed between the MAI and $\delta v$. The correlation coefficient decreases to lower than 0.8 at latitudes higher than $\pm 37.5^\circ$ where the presence of the active region is sparse. From the nature of the two curves, it appears that the linear regression slopes and the magnitude of the correlation are well related. A similar result has been reported for frequency shifts in global modes (Tripathy et al. 2007). For a robust comparison, it will be worth analyzing the entire solar cycle 24 in terms of the ascending and descending phases when the data is available.

### Table 4

| Index | Min | Max | Start | End | Length |
|-------|-----|-----|-------|-----|--------|
| $\delta v$ | 100.2 | 1455.4 | 2007 Jun | 2008 Oct | 17 |
| MAI | 0.8 | 5.0 | 2007 May | 2009 Oct | 30 |
| $F_{10.7}$ | 82.8 | 241.7 | 2007 Feb | 2010 Mar | 38 |
| ISN | 0.0 | 159.5 | 2007 Apr | 2010 Jan | 34 |
| $\delta v$ | 180.8 | 1324.4 | 2007 Jul | 2010 Sep | 39 |
| MAI | 0.9 | 4.8 | 2006 Aug | 2009 Jun | 35 |
| HSN | 0.0 | 80.5 | 2006 Mar | 2009 Sep | 43 |
| $\delta v$ | 28.7 | 1700.0 | 2007 Jun | 2008 Oct | 17 |
| MAI | 0.7 | 5.6 | 2008 Mar | 2010 Feb | 24 |
| HSN | 0.0 | 92.9 | 2007 Sep | 2010 Jun | 34 |

**Note.** Units of $\delta v$ are nHz, MAI is Gauss, and $F_{10.7}$ is sfu. The length is measured in months.

#### 3.4. Minimum Phase

In order to better characterize and quantify the differences between the frequency shifts and activity proxies during the minimum phase, we calculate the length of the minimum period by defining a quiet phase as:

$$(\text{maximum}(x) - \text{minimum}(x)) \times 0.07 + \text{minimum}(x),$$

where $x$ represents either frequency shifts or activity proxies. The factor of 0.07 is chosen based on the observed quiet phase in ISN and $F_{10.7}$. The temporal evolution of $\delta v$, MAI, $F_{10.7}$, and ISN showing the extent of the minimum phase is illustrated in Figure 13 via filled green area. It is evident that the minimum phase in $\delta v$ lasted for a brief period compared to the activity. The minimum and maximum values of these parameters, along with the start time, end time, and the length of the quiet phase are listed in Table 4. The minimum phase in frequency shifts
lasted for a period of about 16 months while all the activity proxies lasted for a period of about 30 months or longer. We have also computed these parameters separately for the northern and southern hemispheres. The values are listed in Table 4. Confirming to known results, the activity (and hence the shifts) are found to be stronger in the southern hemisphere, compared to the northern hemisphere. We also note that the minimum phase lasted longer in northern hemisphere compared to the southern hemisphere, which is probably an indirect manifestation of the stronger activity in the southern hemisphere. We also find an exception in the northern hemisphere where the length of the quiet phase seen in $\delta \nu$ lasted slightly longer compared to the magnetic activity.

We further computed the length of the quiet phase for each latitude. The starting and ending epoch and the length of the minimum phase for $\delta \nu$ and MAI are listed in Table 5; values above $\pm 37.5$ latitudes are omitted since reliable measurements could not be obtained, probably due to small variations between the minimum and maximum phases. Confirming the earlier findings, the span of the quiet phase seen in frequency shifts is found to be reduced, compared to the magnetic activity.

Figure 13: Temporal evolution of (a) $\delta \nu$, (b) MAI, (c) $F_{10.7}$, and (d) ISN showing the extend of the solar minimum phase (filled area), which are overlaid on the running mean over 181 points.
### Table 5

| Latitude | MAI Start | MAI End | Length | δν Start | δν End | δν Length |
|----------|-----------|---------|--------|-----------|---------|-----------|
| 0°:0     | 2008 Jun  | 2011 Apr| 35     | 2009 Mar  | 2011 Dec| 34        |
| 7°:5     | 2006 Apr  | 2010 Jun| 51     | 2006 Dec  | 2010 Sep| 48        |
| 15°:0    | 2006 Feb  | 2009 Sep| 44     | 2006 Apr  | 2009 Sep| 42        |
| 22°:5    | 2006 Jan  | 2009 Jan| 37     | 2006 Feb  | 2008 Jul| 30        |
| 30°:0    | 2006 Mar  | 2008 Sep| 31     | 2005 2008 Mar| 30     |
| 37°:5    | 2007 Jan  | 2009 Sep| 33     | 2005 2008 Jan| 30     |
| 43°:5    | 2008 Aug  | 2011 Jun| 35     | 2009 2011 Sep| 33     |
| 50°:0    | 2008 Mar  | 2010 Jul| 29     | 2008 2010 Jul| 29     |
| 57°:5    | 2006 Apr  | 2009 Oct| 41     | 2006 2010 Jan| 40     |
| 64°:5    | 2006 Apr  | 2009 Apr| 37     | 2006 2008 Dec| 36     |
| 71°:5    | 2007 Mar  | 2009 May| 27     | 2007 2008 Dec| 24     |

**Note.**

* Length is measured in months.

### 4. CONCLUSIONS

In this work we have analyzed high-degree mode frequencies of the Sun over a period spanning about 13 years to understand their behavior with the solar cycle. These frequencies are obtained from the application of the ring-diagram technique to the high-resolution Doppler observations made by the GONG and cover the period of 2001 July 26 to 2014 June 19.

We compared the mean frequency shifts with three activity proxies — the 10.7 cm radio flux, the International sunspot Number, and the strength of the local magnetic field measured by the MAI. In general, a good agreement is found between the shifts and activity indices and the correlation coefficients are found to be comparable with intermediate-degree mode frequencies. Although we do not have data covering all of cycle 23 or cycle 24, the phase-wise variations of the shifts indicated that the frequencies responded to solar activity equally well in both the cycles, although the shifts per unit activity hinted that cycle 24 is weaker than cycle 23. We also noticed that the frequency shifts do not show a prolonged minimum phase as is seen in many activity indicators, suggesting weak correlations between the shifts and activity indices during the activity minimum period.

For the first time, we also computed frequency shifts over the north and south hemispheres separately and analyzed them as a function of HSN and MAI. The asymmetric nature of the magnetic activity is also reflected in the frequency shifts. We note that the shifts behaved similarly in the northern and southern hemispheres until the end of 2007 and diverged thereafter. The Sun had started showing deviations even before the onset of the minimum phase and that the mode frequencies sensed these changes before they otherwise manifested in the visible surface. Furthermore, we find that the shifts reached their maximum amplitude at two different periods—the amplitude of shifts in the northern hemisphere peaked during late 2011, more than two years earlier than in the south. Such asymmetric behavior between the ascending and descending phases of solar cycle 23 has been studied by Simoniello et al. (2013a) and it has been argued that an increase of asymmetry might be used as a precursor of a grand minima. Theoretical calculations also indicate that strong asymmetries are induced by the quadrupolar component and characterize periods before entering or leaving the grand minima state (Tobias 1996; Knobloch et al. 1998). Thus, in our view the asymmetry associated with the frequency shifts may have certain implications for the dynamo action taking place either in the sub-surface shear layer or at the base of the convection zone.

Since the shifts and the hemispheric activity indices are found to be significantly correlated, we suggest that the frequency shifts may be used as an indicator of hemispheric activity since not many indices are measured over the two hemispheres separately. We also investigate the variations at different latitudes and conclude that the shifts in active latitudes correlate well with the local MAI. We believe that this analysis provides an impetus to analyze solar cycle 24 in terms of hemispheric asymmetry and different phases in order to comprehend the magnetic cycle of the Sun.

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