The natural science of cosmology

P J E Peebles

Joseph Henry Laboratories, Princeton University, Princeton NJ 08544, USA
E-mail: PJEP@Princeton.edu

Abstract. The network of cosmological tests is tight enough now to show that the relativistic
Big Bang cosmology is a good approximation to what happened as the universe expanded and
cooled through light element production and evolved to the present. We have explained the
reasons to reach this conclusion, commented on the varieties of philosophies informing searches
for a still better cosmology, and offered an example for further study, the curious tendency of
some classes of galaxies to behave as island universes.

1. Introduction
Cosmology aims to draw large conclusions from exceedingly limited evidence, but we have to
conclude, to our surprise, that the considerations reviewed in Section 2 make a convincing
case that the relativistic Friedman-Lemaître expanding world model is a good approximation.
Cosmology is a work in progress, of course, and many lines of research are aimed at establishing
a still better theory. We have offered in Section 3, thoughts on directions this research has taken
and likely will take. One direction is the search for anomalies, where the reconciliation of theory
and observation seems particularly difficult. A good anomaly will teach us something, either
a better understanding of the standard model or clues to an improved one. In Section 4, we
have given an example, for the apparent behaviour of pure disk and elliptical galaxies as island
universes.

2. The cosmological tests
The standard ΛCDM cosmology assumes the general theory of relativity. This is an extrapolation
of some 14 orders of magnitude in length scale from the precision tests on the scales of the
Solar System and smaller. It assumes that 95% of the present mass of the universe is in two
hypothetical forms, dark matter and dark energy. And it assumes initial conditions that agree
with simple versions of the inflation picture of the very early universe, which is encouraging, but
inflation is an idea, not a theory. Substantiating this considerable variety of bold assumptions
is a demanding challenge. But the considerations reviewed here, lead us to conclude that the
challenge has been met, largely.

Detecting the relativistic curvature of the redshift-magnitude relation [1, 2] is a rightly
celebrated contribution to the cosmological tests. But one gets some feeling for the challenge
of substantiating cosmology by imagining the situation in an alternative universe, in which the
only significant cosmological test is the redshift-magnitude relation. That would include the
demonstration of a close to linear relation between redshift and distance at low redshift, which
is in line with Einstein’s argument for homogeneity. We have expected the demonstration of the
departure from linearity at high redshift by the SNe Ia measurements would split the cosmology community in this alternative universe into three camps. Those who already favoured the Steady State cosmology on philosophical grounds would point out that the curvature is in line with the Steady State prediction, and would rightly claim strong support for this cosmology. Among those who had a philosophical preference for the relativistic Friedman-Lemaître cosmology one camp would add Einstein’s cosmological constant \( \Lambda \) to their cosmology. The other camp would point out that the indicated value of \( \Lambda \) is absurdly different from any reasonable estimate from well-established quantum physics. They would point out that in the Friedman-Lemaître cosmology, the universe at \( z > 1 \) was different from now, and SNe Ia then, therefore, had to have been at least somewhat different from present-day ones. They would acknowledge that carefully explored tests have given no indication of evolution of the properties of SNe Ia, but still insist that supernovae cannot be examined in close detail, and there is no way to exclude the possibility that some subtle evolution has escaped the tests. Each of these three camps has a reasonable case.

**Figure 1.** Measurements of the CMB spectrum, compiled by Alan Kogut. The dotted curve is a thermal Planck spectrum.

The point is that establishing cosmology requires a network of tests that check for systematic errors, both in the measurements and in the theory used to interpret the measurements. For example, suppose cosmologists in the alternate universe had, in addition to the redshift-magnitude relation, measurements of the cosmic microwave background radiation (the CMB), which show that the CMB is isotropic to better than a part in \( 10^5 \) and that the spectrum is very close to thermal, as shown in Figure 1 (compiled and plotted by Alan Kogut [3], NASA GSFC). The only reasonable way to reconcile this measured thermal spectrum with the classical Steady State cosmology is to postulate that the universe is optically thick to the absorption and emission of microwave radiation at very small redshifts, for otherwise the CMB spectrum would be an unacceptable mix of temperatures set by the distribution of redshifts summed out past
unit optical depth. Since people in the alternate universe could measure the CMB, they would have detected discrete sources at CMB wavelengths. If they had adequate optical telescopes, they would see that some were distant radio galaxies, meaning the universe was optically thin. This serious challenge to the classical Steady State cosmology is an example of a contribution to a network of tests which, if dense enough, may establish cosmology.

Returning to our universe, let us note that no natural science is complete, and that certainly includes the established relativistic cosmology. For example, we have a few constraints on what happened in the very early universe. One can save the Steady State philosophy by postulating that our universe behaves in a Quasi-Steady State manner [4], where the continual creation of the classical version is replaced by periodic bursts of creation at densities large enough to be capable of thermalizing the radiation. If the bursts of creation built up an appropriate distribution of matter and radiation at densities greater than that of the Big Bang at light element nucleosynthesis, and were followed by the expansion described by the relativistic model, then QSS would pass all cosmological tests, and merit a place with variants of inflation in models for the very early universe that were viable but need work.

Because the value of Einstein’s Λ in the standard Big Bang cosmology seems so unlikely within quantum physics, it is natural and important to consider whether the evidence for detection of Λ results from systematic error in the application of the relativistic cosmology. An example is the proposal that the strongly inhomogeneous distribution of mass on scales less than a few tens of megaparsecs may significantly affect the general rate of expansion of the universe [5]. We believe, Siegel and Fry [6] have first answered this question. One may start with the observation that deflection angles δθ, in gravitational lensing are small, δθ ∼ GδM/rc² ≲ 10⁻⁵ around a mass concentration δM of radius r. This angular deflection is a measure of the departure from the homogeneous Friedman-Lemaître geometry |δgμν| ≲ 10⁻⁵gμν. (Black holes are exceptions but irrelevant, because at any distance of interest for this discussion they behave as ordinary particles.) One is invited to compute the spacetime geometry to second order in δgμν sourced by the mass auto-correlation function, which is reliably measured on the scales of significantly non-linear mass clustering. The result (which requires a modest extension of the calculation in [6]) is that non-linear mass fluctuations modelled as a gas of particles with mass m and mean number density n̅ appear in the Friedman equations as an effective contribution to the mean mass density ρeff = m̅n̅(1 + K + W), and to the pressure peff = m̅n̅(2K + W)/3, where K = (ν²)/2 is the kinetic energy of peculiar motion per unit mass and W = −1/2Gm̅n̅a² ∫ d³x⟨δ(r)δ(r + x)⟩/x is the gravitational energy per unit mass belonging to the departure from homogeneity. The effective pressure peff can be negative as for Λ, but it is tiny, five orders of magnitude smaller than the standard estimate of Λ. This is not a way out of Λ or dark energy.

Another proposed way out supposes that Λ is small or zero, and that the mass density and expansion rate vary with distance from us in such a way as to make the angular size distance vary with redshift in the same way as in a homogeneous universe with the standard value of Λ. We would have to be very near the centre of an accurately spherically symmetric universe, for otherwise the Sachs-Wolfe effect would produce an unacceptably large CMB anisotropy. Our favoured position is arguably unlikely, but more important is that this model is challenged by other cosmological tests. Intergalactic plasma scatters CMB photons, meaning the spectrum in Figure 1 is a mix of spectra originating at a range of distances. If the temperature at decoupling were not very close to homogeneous, the CMB spectrum would be a mix of thermal spectra at different temperatures, contrary to the measurements. In the standard cosmology, acoustic waves in the baryons and radiation prior to decoupling produced waves in the CMB anisotropy power spectrum, and in the power spectrum of the large-scale galaxy distribution. Analyses [7, 8] pointing to consistency of the measurements of the CMB and galaxy distributions depend on, and test, many elements of the cosmology, but our point is that measurements of the angular distribution of the radiation probe what happened near the edge of the observable universe, and
measurements of the spatial distribution of the galaxies probe what happened much closer to us. The consistency of these measurements within the homogeneous ΛCDM cosmology argues that conditions actually were quite similar across the sky at these two quite different ranges of distance, or else that there was yet another conspiracy.

Figure 2. Figures taken from papers on cosmological tests, for the purpose of illustrating the broad variety of phenomena, and means of measurement and analysis in the network of tests.

This discussion could continue at length, but we, instead, refer to Figure 2 for a schematic illustration of the broad variety of phenomena and means of measurement employed in the cosmological tests. The consistency of results from these many and varied approaches convincingly argues against the possibility that we have been seriously misled by systematic errors in the observations or errors in the theory by which the observations are interpreted.

We hope it is understood that we are not arguing that cosmology after baryogenesis is complete: that is, contrary to the experience that natural sciences advance by successive approximations. But we can be sure that if there is a better cosmology, it will predict a universe that looks much like ΛCDM, because the tests look at the universe from many sides now, and find that it looks much like ΛCDM.
3. The future
We have got some sense of the future of this subject by recalling how we got to where we are now. We have seen several paths common to cosmology and other parts of natural science.

3.1. Serendipitous discovery of suggestive phenomena
In relativity physics, one might think of the apparently curious behaviour of the velocity of light. In cosmology, one thinks of the linear redshift-distance relation, which by 1936 Hubble and Humason [9] had tested out to \( z = 0.1 \). This relation was the centrepiece of most discussions of cosmology until the CMB, in 1965. Serendipity strikes seldom, but let us remember that it happens.

3.2. The philosophical appeal of ideas
Einstein followed this path to wonderfully good effect in his discovery of general theory of relativity, and his proposal that a philosophically satisfactory universe is homogeneous in the large-scale average. He had no justification for homogeneity from phenomenology, and only a few hints to GR: Electromagnetism and the Eötvös experiment. It is striking that Nature on occasion agrees with our ideas of elegance. A contrary example is the fate of the classical Steady State cosmology, though the idea arguably reappears in eternal inflation. We have included as modern contenders cosmic strings, superstring cosmology, the multiverse and the anthropic principle.

3.3. Mathematical incompleteness
We take it that Nature abhors the singularities in general relativity. The community favourite remedy for the singularity in ΛCDM at infinite redshift is a variant of inflation, which needs work.

3.4. Testing ideas
The central example in cosmology is the programme of cosmological tests that has succeeded remarkably well, after some re-negotiation of ideas about Λ. An influential idea driving research now is that there may be a signature of inflation in the polarization of the CMB. A successful detection would have a wonderfully strong effect on the status of inflation. Laboratory searches may detect dark matter particles that behave much like the hypothetical cold dark matter in ΛCDM. If this detected matter along with massive neutrinos saturates the mass in matter in the dark sector, this will close a line of research as far as cosmology is concerned. But resolution of an issue more often raises new ones. May be some component of the dark matter will prove to do something interesting, leading us to new lines of research. The fascinating enigma of Einstein’s Λ motivates the study of modifications of the physics of gravity and the dark sector, and the research programmes aimed at detecting or bounding evolution of Λ. Testing variation of Λ is good science, though we must say, we are uneasy about such heavy expenditure of resources on a shot in the dark. Let us support it, but also support wider casts of the net.

3.5. Anomalies
Two celebrated examples are in Lord Kelvin’s essay [10], *Nineteenth Century Clouds over the Dynamical Theory of Heat and Light*. One cloud is exemplified by the curious behaviour of mercury vapour. A spark causes the vapour to radiate at sharply defined frequencies seemingly scattered at random across the visible spectrum. The vapour must contain resonators that radiate at definite frequencies when excited. The anomaly is that these radiators ought to be excited when the vapour is heated, yet the heat capacity of mercury vapour is quite close to that of an ideal gas of structureless particles. Lord Kelvin knew physics is well founded: He owed
his fortune and peerage to his great contributions to physics, to say nothing of his income from physics-based patents for transatlantic telegraph cables among other things. Our impression of Kelvin’s position, put in modern jargon, is that physics is solid and the rates for thermal excitation of internal energies of gasses must be slow enough that heat capacities are measured in systems that have not relaxed to statistical equilibrium. At the time Planck was struggling with another heat capacity problem, that of thermal radiation, another precursor of the perfect storm of new physics to come. Kelvin acknowledged that the other cloud in his essay, the curious behaviour of the velocity of light, is “very dense”. He mentioned with approval “the brilliant suggestion made independently by Fitz Gerald and by Lorentz of Leiden that the motion of ether through matter may slightly alter its linear dimensions.” That was yet another precursor of new physics, though we imagine Kelvin had in mind a mechanical effect of the ether on the matter, to be dealt with in the standard physics of the day.

We put Milgrom’s [11] MOND (Modified Newtonian Dynamics) in this path. His idea that flat rotation curves of late-type galaxies could signify a departure from Newtonian gravity rather than the presence of dark matter seems quite out of line with the cosmological tests discussed in the last section, but MOND gives a remarkably successful prediction of the relation between circular velocity and detected baryonic mass in stars and neutral gas in these galaxies [12]. One has to be impressed that this prediction was observed over a far broader range of masses than was known when Milgrom proposed MOND. We should support continued research on this cloud over cosmology, by a few capable people, for it may teach us something of value. We offer in the next section two other anomalies that, whether real or apparent, also seem likely to be informative.

3.6. Learning to compute
Consider Kelvin’s opinion of heat capacities, and the un-enthusiastic community response to MOND. We have a well-founded cosmology, just as Kelvin had a well-founded physics. Kelvin saw the need for better computation of rates of relaxation to statistical equilibrium of internal energies of gasses. The cosmology community by and large is optimistic that better methods of computation will show that the standard cosmology predicts the observed baryon mass-circular velocity correlation, and the curious behaviour of the dwarf galaxy spatial distribution and baryonic mass function [13], and other such issues. This echoes the old admonition in physics, “just trust the theory”, which often proves to be a useful guide. Trust in the standard cosmology, as in explaining the MOND prediction within ΛCDM, is a good idea too, but for a different reason. The cosmological tests are far less dense than tests of the more well developed branches of natural science, and trust, therefore, more questionable. But by concentrating almost exclusively on the one standard cosmology rather than scattering attention across alternatives, the community has a far better chance of discovering whether within the available evidence there really is something wrong with the chosen standard model. On the other hand, we should bear in mind the outcome of more flexible thinking about Kelvin’s clouds. It is good strategy for some to scout more flexible thinking about cosmology, which may aid discovery of problems with given ideas, if there are any, and advance discovery of remedies, if they are needed.

3.7. Adding decimal places
In Lord Kelvin’s time, some felt that all that was left to physics was adding decimal places. But this example reminds us why improving accuracy with no other goal in mind is not to be deprecated: it can reveal anomalies. And despite the optimistic new title for our subject, precision cosmology, we still are proud to see measurements that get beyond the first decimal place.
4. Island universes
We have offered two examples along the path mentioned in Subsection 3.5, apparent anomalies in the behaviour of pure disk galaxies and elliptical galaxies. We have mentioned these particular cases because both seem to us to be resistant enough to explanation within the ΛCDM cosmology to be very likely to teach us something.

A pure disk galaxy is shown in Figure 3. The main image of this nearby face-on spiral galaxy was made by R. Gendler by piecing together Hubble Space Telescope images and ground-based data. The dark lanes in the spiral arms are caused by obscuration by interstellar dust. The inset is a single HST image obtained by John Kormendy et al [14]. It shows that the dust lanes run inward to a central star cluster with radius $\sim 0.1$ per cent of the radius of the luminous part of the galaxy. This galaxy appears to be in a near relaxed state, in which the dust will have settled onto the plane of the disk. The prominent dust lanes, extending very close to the centre of the galaxy indicate the bulk of the stars, also are confined to the plane, and that the structure of this galaxy is largely supported by rotation. Kormendy et al have argued that this situation does not seem to be uncommon, because they have found that about half the 20 nearest big galaxies are largely supported by motions in the plane. The striking whirlpool appearance of M 101 seems to require that this galaxy grew by dissipative settling of gas or plasma onto the growing disk prior to significant star formation. If stars had formed in significant amounts in matter before it had settled into the disk it would have populated a stellar halo or classical bulge, which were not prominent components of this or other pure disk galaxies.

Numerical simulations of structure formation in ΛCDM seem to predict that galaxies were assembled by hierarchical mergers of fragments (as illustrated in Figure 3 in [15]). The larger fragments would have been massive and dense enough that pressure support required matter temperatures large enough to have collisionally ionized hydrogen. The plasma would lose energy by thermal bremsstrahlung, causing the fragments to contract until some energy source, presumably star formation, prevented it. But, if too many stars formed in these fragments, it would have produced a distinct stellar component outside the disk in the final product of the hierarchy of mergers. That is quite acceptable for some galaxies, such as our nearest large neighbour M 31, which has the substantial bulge of old stars that one might have expected to
have resulted from the merging seen in structure formation simulations. Recent simulations [16, 17] do produce impressive pure disk galaxies, through careful attention to models for the conditions that promote and prevent star formation, but the mystery remains. The global star formation rate density was at its peak at $4 \gtrsim z \gtrsim 2$. How could the progenitor fragments of pure disk galaxies have “known” not to have participated in this generally high global star formation rate? One piece of the matter tumbling together according to the ΛCDM picture of the formation of the pure disk galaxy in Figure 3 “knew” it was going to host the growing disk, and start growing it at redshift well above unity if the age of the disk of the Milky Way [18] was typical of pure disk galaxies, while the rest of the fragments “knew” they had to hold off star formation until they had reached the growing disk. It is a curious situation.

![Figure 4. Comparison of relations among elliptical galaxies found in regions of lower and higher ambient densities.](image)

A possibly related issue is why overall properties of another large class of galaxies, ellipticals, are so little sensitive to environment. The phenomenon is illustrated in Figure 4, taken from Nair, van den Bergh and Abraham [19] (and based on the Sloan Digital Sky Survey). The vertical axis is a measure of radius. The authors have taken special care with this parameter, because galaxies do not have edges: they just trail off into intergalactic space. The sample is separated according to ambient density, lower in the left-hand panel, higher in the right-hand panel. One sees several classical regularities. The Kormendy relation is the tendency for ellipticals with larger luminosities $L$ to have larger radii, as indicated by the solid lines. The Faber-Jackson relation is the tendency for ellipticals with larger stellar velocity dispersions $\sigma$, as one can make out from the colour coding for $\sigma$. The figure also illustrates the preference of more luminous ellipticals for denser environments: In the higher ambient density sample in the right-hand panel, the data points more densely populate the scatter plot near luminosity $\sim 10^{11}$ Solar luminosities and extend beyond that to luminosities three times the largest in the lower ambient density sub-sample. This preference is predicted by the ΛCDM structure formation simulations. Also, Nair et al [19] have found that radius $R_e$ that contains half the starlight of an elliptical shows more scatter, and more distinctly so in the lower density subset, than the radius $R_{90}$ used in Figure 4 that contains a larger fraction of the starlight. This is an environmental effect, perhaps rearrangement of internal structure by galaxy interactions [19]. Environment certainly matters, but to us the striking and surely significant phenomenon is the overall insensitivity of the elliptical galaxy size-luminosity relation to the present environment.

We have seen a similar phenomenon in the Zhu, Blanton and Moustakas [20] comparison of mean spectra of central parts of ellipticals classified by the central stellar velocity dispersion $\sigma$.
Figure 5. Comparison of mean spectra of elliptical galaxies found in regions of lower and higher ambient densities for three ranges of stellar velocity dispersion.

and ambient density. Figure 5 (kindly provided by Gunagtun Zhu using the results in [20], and based on SDSS) shows ratios of mean spectra in the three bins in \( \sigma \) indicated in the two panels, separately for the higher and lower density sub-samples. The ratios of spectra, with the lower \( \sigma \) divided by the larger \( \sigma \) mean spectrum, increase to the left, toward the blue. This illustrates the classical colour-magnitude relation: More luminous ellipticals, that tend to have larger \( \sigma \), tend to be redder. The scale is larger in the lower panel that compares the lowest to highest \( \sigma \) bins than in the upper panel that compares the intermediate to highest \( \sigma \) bins; the relation is monotonic across the sample. The ratios for the lower density (field) and higher density (rich group) subsets are plotted in blue and red. The blue spike in the lower panel is a hydrogen recombination line likely from star-forming regions that are more common in ellipticals in lower density regions. The wonderful SDSS stability also reveals in these data a very clear relation between spectrum and ambient density, as shown by the systematic tilt of blue relative to red curves. Figure 11 in [20] shows that at given \( \sigma \) the spectra are bluer at lower density. There are exceptions to the very modest sensitivity to environment illustrated in Figure 5. The brightest galaxy in a cluster, which usually is near the cluster centre defined by the X-ray emission [21], tends to be redder than other early-type cluster members with the same \( \sigma \) [22], or it may host significant star formation [23]. That is, environment certainly matters to ellipticals. But again, to us the central point is that in the general class of ellipticals, the variation of spectrum with environment is wonderfully small compared to the variation with velocity dispersion.

We may say that Nair et al have classified ellipticals by a measure of internal conditions, the luminosity \( L \), and a measure of external conditions, the ambient density. They have found that another measure of internal conditions, the radius, is much more sensitive to \( L \) than to the
ambient density. Zhu et al have classified by the measure $\sigma$ of internal conditions and a measure of external conditions, and found that another measure of internal conditions, the spectrum, is much more sensitive to $\sigma$ than ambient density. Other examples along the same lines are the relation between colour and luminosity [24], between red galaxy mass-to-light ratio and radius [25, 26], and between galaxy stellar mass and radius [27]. One might read this to mean that elliptical galaxy evolution has an attractor behaviour that pulls ellipticals with given present $L$ to near common spectra, radii, and mass-to-light ratios. This is conceivable, and good science requires us to keep it in mind, but it seems unlikely to the author. The alternative is that the present environment does not much matter to ellipticals either, because they have been formed at high redshift, when there is little variation in environment, and then passively evolved with little interaction with the environment, or else they grow more slowly, but by rearrangement of matter that is already part of the proto-galaxy evolving as a closed system. This straightforward reading of the evidence is not what pure dark matter simulations have led us to expect from $\Lambda$CDM.

Other broad classes of galaxies do show signatures of interaction with their environment. The transition of spiral galaxies into more nearly gas-free S0 galaxies seems to be a particularly clear example [28, 29]. Galaxies are observed in the act of merging, and theory and intuition argue that the numerous close galaxies at high redshift will have suffered many more mergers [30]. These observations are in line with the hierarchical assembly of structures by mergers seen in simulations of the standard cosmology, and it is entirely reasonable, therefore, that mergers have figured heavily in discussions of how the galaxies formed. Less reasonable, we think, is the scant attention to the evidence that two broad classes of galaxies, pure disks and ellipticals, have evolved in near isolation from their surroundings, as island universes. This aspect of galaxy evolution also must be telling us something important about the origin of cosmic structure.

5. Concluding remarks
Cosmology has advanced far more than we, or suppose anyone, imagined a half century ago when we started looking into this subject. Perhaps witnessing the many wrong turns taken, and right turns overlooked so long, has conditioned us to be more doubtful about the stability of the present standard $\Lambda$CDM cosmology than we have seen in the thinking of many younger colleagues. But we can agree that the advances of the observational and theoretical bases are substantial and of lasting importance. This review leaves us with the impression that there is no shortage of interesting problems for study by the next generation. We have mentioned the island universe puzzle, because it seems to us that the opportunities for research are even broader than generally advertised.

Acknowledgements
The author has benefitted from discussions with Michael Blanton, John Kormendy, Yen-Ting Lin, Piero Madau, Chris McKee, Surlud More, John Moustakas and Gunagtun Ben Zhu.

References
[1] Riess A G et al 1998 Astron. J. 116 1009
[2] Perlmutter S et al 1999 Ap. J. 517 565
[3] Kogut A J 2011 Private communication
[4] Narlikar J V 1997 J. Ap. & A. 18 353
[5] Kolb E W, Matarrese S and Riotto A 2006 New J. Phys. 8 12
[6] Siegel E R and Fry J N 2005 Ap. J. Lett. 268 L1
[7] Komatsu E et al 2011 Ap. J. Supp. Ser. 192 18
[8] Percival W J et al 2001 MNRAS 327 1297
[9] Hubble E 1936 The Realm of the Nebulae (New Haven: Yale University Press)
[10] Kelvin (Thomson W) 1901 Phil. Mag., Sixth Ser. 2 1
11

[11] Milgrom M 1983 *Ap. J.* 270 365
[12] McGaugh S S 2012 *Astron. J.* 143 40
[13] Macciò A V, Kang Xi, Fontanot S, Somerville R S, Koposov S and Monaco P 2010 *MNRAS* 402 1995
[14] Kormendy J, Drory N, Bender R and Cornell ME 2010 *Ap. J.* 723 54
[15] Peebles P J E and Nusser A 2010 *Nature* 465 565
[16] Brook C B *et al* 2011 *MNRAS* 415 1051
[17] Guedes J, Callegari S, Madau P and Mayer L 2011 *Ap. J.* 742 76
[18] Aumer M and Binney J J 2009 *MNRAS* 397 1286
[19] Nair P, van den Bergh S and Abraham RG 2011 *Ap. J. Lett.* 734 L31
[20] Zhu G, Blanton M R and Moustakas J 2010 *Ap. J.* 722 491
[21] Lin Y-T and Mohr J J 2004 *Ap. J.* 617 879
[22] Roche N, Bernardi M and Hyde J 2010 *MNRAS* 407 1231
[23] O’Dea K P *et al* 2010 *Ap. J.* 719 1619
[24] Hogg D W *et al* 2004 *Ap. J. Lett.* 601 L29
[25] Bernardi M, Nichol R C, Sheth R K, Miller C J and Brinkmann J 2006 *Astron. J.* 131 1288
[26] Magoulas C, Colless M, Jones H, Mould J and Springob C 2010 *Highlights Astron.* 15 84
[27] Maltby D T *et al* 2010 *MNRAS* 402 282
[28] Just D W, Zaritsky D, Sand D J, Desai V and Rudnick G 2010 *Ap. J.* 711 192
[29] Kormendy J and Bender R 2010 *Ap. J. Supp. Ser.* 198 2
[30] López-Sanjuan C *et al* 2012 [arXiv:1202.4674]