Transversity Parton Distribution

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Transversity distribution is one of the three fundamental parton distributions that completely describe polarized spin 1/2 nucleon. Its chiral odd nature prevented for many years its experimental exploration, however presently we have obtained great deal of information about this distribution. This includes experimental data from Semi Inclusive Deep Inelastic Scattering, knowledge of scale dependence and phenomenological extractions. I will discuss main features of this distribution and indicate the future improvements of our knowledge.

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1. Introduction

Distribution of partons in a polarized spin-1/2 hadron can be completely described by three collinear Parton Distribution Functions (PDFs): unpolarised parton distribution $f_1$, helicity distribution $g_1$, and transversity distribution $h_1$. These standard collinear PDFs are defined through collinear factorization theorems and obey DGLAP evolution equations [1, 2, 3, 4, 5, 6].

The transversity distribution [7] describes transversely polarized quarks in a transversely polarized nucleon. Formally such a distribution can be expressed by

$$h_1^q(x) = \int \frac{d\xi}{2\pi e^{-ixP + \xi S_T}} \mathcal{M}(P, S_T)$$

(1.1)

The corresponding charge is called “tensor charge”

$$\delta q = \int_0^1 dx (h_1^q(x) - h_1^\bar{q}(x))$$

(1.2)

Transversity obeys so-called Soffer bound [8]

$$|h_1(x, Q^2)| \leq \frac{1}{2} (f_1(x, Q^2) + g_1(x, Q^2))$$

(1.3)

This bound was shown to be preserved at LO accuracy in Ref. [4] and at NLO accuracy in Ref. [5].

Transversity is the least known of the three collinear distributions and the reason is that it cannot be measured in Deep Inelastic Scattering (DIS) due to its chiral odd nature. It should couple to another chiral odd quantity (chiral odd fragmentation or chiral odd distribution function, for example transversity itself). The best channel to measure transversity remains polarized Drell-Yan (preferable proton anti proton) process in which one could measure the product of transversity distributions directly, see Ref. [9]. QCD evolution of collinear transversity distribution is well known, see Refs. [4, 5, 6]. It does not couple to gluons and thus exhibits non-singlet $Q^2$ evolution. Gluon transversity distribution does not exist either. This leads to the fact that transversity is suppressed at low-$x$ makes it a valid object to study in high-$x$ region by Jefferson Lab 12 [10].

Currently the knowledge on transversity comes from Semi Inclusive Deep Inelastic Scattering (SIDIS) experimentally observed at HERMES [11, 12], COMPASS [13, 14] and JLab 6 [15] in single spin asymmetries where transversity couples to so-called Collins fragmentation function [16]. Information on the convolution of two chiral-odd fragmentation functions is obtained from $e^+e^- \rightarrow h_1 h_2 X$ processes [17, 18, 19]. One usually measures low transverse momentum final hadron and thus one applies Transverse Momentum Dependent factorization. The transversity in this case depends also on intrinsic transverse motion of quarks $k_{\perp}$ and one speaks of Transverse Momentum Dependent (TMD) transversity. One can also study transversity coupled to so-called di-hadron fragmentation function [20, 21, 22].

The $u$ and $d$ quark transversity distributions, together with the Collins fragmentation functions, have been extracted for the first time in Refs. [23, 24], from a combined analysis of SIDIS and $e^+e^-$ data. The most recent extraction is presented in Ref. [25].

Di-hadron method was implemented in the analysis of Ref. [26] and the results on the extraction of transversity from Refs. [23, 24, 25] and Ref. [26] agree with each other quite well.

QCD evolution of Transverse Momentum Dependent transversity was recently obtained in Ref. [27].
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2. Phenomenology

The result on extraction of the transversity is presented in Fig. 1.\footnote{The plot is from Ref. \cite{25}} One can see that $u$ quark transversity is positive and $d$ quark transversity is negative. This result is coming from global analysis of SIDIS HERMES \cite{11,12}, COMPASS \cite{13,14} and $e^+e^-$ BELLE \cite{17,18,19} data.

Experimentally so-called Collins asymmetry in SIDIS with unpolarised beams ($U$) and transversely polarized target ($T$) is measured and it is proportional to convolution of transversity and Collins fragmentation functions

\[ A_{UT}^{\sin(\phi_h + \phi_S)} \propto \sum_q h_T^q \otimes H_{1q} \tag{2.1} \]

here $\phi_h$ and $\phi_S$ are azimuthal angles of produced pion and polarization vector, experimentally observed modulation is proportional to $\sin(\phi_h + \phi_S)$ and sign $\otimes$ denotes usual TMD convolution \cite{28}.

One can see that knowledge of Collins fragmentation function ($H_{1q}^\perp$) is needed in order to extract transversity, fortunately in $e^+e^-$ process one observes an asymmetry which is related to convolution of two Collins functions $\sum_q H_{1q}^\perp \otimes H_{1\bar{q}}^\perp$. This allows us to have perform global analysis \cite{23,24,25} of SIDIS and $e^+e^-$ data.

We also present the results on the tensor charge at $Q^2 = 0.8 \text{ GeV}^2$ in Fig. 2.\footnote{The plot is from Ref. \cite{25}}

Future measurements at Jefferson Lab 12 are going to be very important for the extraction of the transversity and the tensor charge. We estimate that corresponding improvement of the statistical error of the extraction will be a factor of 5 approximately \cite{10}. This will mean that from almost 50% uncertainty we will be able to extract tensor charge with 10% uncertainty and compare it better to model predictions.
Figure 2: The tensor charge for \( u \) (left) and \( d \) (right) quarks, computed using the transversity distributions obtained in Ref. [25]. The gray areas correspond to the statistical uncertainty bands of the extraction. The results are compared with those given in Ref. [24] (number 2) of the model calculations of Refs. [29, 30, 31, 32, 33, 34, 35] (number 3-9) and with the results extraction from Ref. [26] (number 10).

Tensor charge is important for some dynamical effects of new heavy Beyond Standard Model degrees of freedom [36]. In order to constrain possible parameters of those models one needs precise knowledge of the tensor charge.

Future Electron Ion Collider [37] will also allow us to study carefully \( Q^2 \) evolution of transversity and explore low-\( x \) region.

3. Conclusions

Interested reader is referred to several reviews that describe the transversity in greater detail, see Refs [38, 39, 40]. I have not discussed all possible ways to access transversity, for instance \( \Lambda \) electroproduction in SIDIS. In this case one needs to know the chiral odd fragmentation function of \( \Lambda \) production and it can be accessed via \( e^+e^- \rightarrow \bar{\Lambda}X \). One could also study transversity in proton proton scattering by utilizing \( pp \uparrow \rightarrow \pi jetX \) [41].

In future we will have data from BABAR Collaboration, which have performed an independent new analysis of \( e^+e^- \rightarrow h_1h_2X \) data [42], Jefferson Lab 12 will provide precision data in high-\( x \) region [10] and thus complement results obtained in SIDIS at HERMES [11, 12], COMPASS [13, 14] and JLab 6 [15].

Generally proton proton scattering can be described by so-called twist-3 factorization in which one studies multi-parton correlations of partons and corresponding Efremov-Teryaev-Qiu-Sterman functions [43, 44, 45, 46, 47, 48]. These functions are related to TMD functions and more globally twist-3 formalism and TMD formalism are closely related to each other, and have been shown to be equivalent in the overlap region where both can apply [49, 50, 51].
Once a comprehensive global analysis of the data from SIDIS (TMD) and proton proton scattering (twist-3) is done (see preliminary results in Refs. [52, 53]) we are going to obtain a complete description of asymmetries and corresponding parton distributions including transversity parton distribution.

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