Chandrayaan-2 Dual-frequency Synthetic Aperture Radar (DFSAR): Performance Characterization and Initial Results

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Abstract

The Dual-Frequency synthetic aperture radar (DFSAR) system manifested on the Chandrayaan-2 spacecraft represents a significant step forward in radar exploration of solid solar system objects. It combines SAR at two wavelengths (L and S bands) and multiple resolutions with several polarimetric modes in one lightweight (~20 kg) package. The resulting data from DFSAR support the calculation of the 2 × 2 complex scattering matrix for each resolution cell, which enables lunar near-surface characterization in terms of radar polarization properties at different wavelengths and incidence angles. In this paper, we report on the calibration and preliminary performance characterization of DFSAR data based on the analysis of a sample set of crater regions on the Moon. Our calibration analysis provided a means to compare on-orbit performance with prelaunch measurements, and the results matched with the prelaunch expected values. Our initial results show that craters in both permanently shadowed regions (PSRs) and non-PSRs that are classified as circular polarization ratio-anomalous in previous S-band radar analyses appear anomalous at the L band also. We also observe that material evolution and physical properties at their interior and proximal ejecta are decoupled. For the Byrgius C crater region, we compare our analysis of dual-frequency radar data with the predicted behaviors of theoretical scattering models. If crater age estimates are available, a comparison of their radar polarization properties at multiple wavelengths similar to that of the three unnamed south polar crater regions shown in this study may provide new insights into how the rockiness of craters evolves with time.

Unified Astronomy Thesaurus concepts: The Moon (1692); Lunar craters (949); Polarimetry (1278); Radar astronomy (1329); Lunar science (972); Lunar features (953)

1. Introduction

The physical properties of lunar surface materials provide an independent means to help unravel lunar geologic history and evaluate their exploration potential. Additionally, near-surface physical properties of the regolith are of importance to address significant questions in lunar geology, such as understanding the spatial limits and relative timing of mare flow units and impact-related processes that eventually influence spatial differences in surface and subsurface composition (e.g., Cahill et al. 2014). In this context, polarimetric radar backscatter data for the Moon provides valuable insights regarding the physical properties of the near surface with particular sensitivity to ilmenite content, and surface or buried rocks with a diameter of about 1/10 the radar wavelength and larger (e.g., Carter et al. 2009; Campbell et al. 2010; Campbell 2012).

The Dual-Frequency Synthetic Aperture Radar (DFSAR) aboard India’s second lunar mission Chandrayaan-2 is the first fully polarimetric synthetic aperture radar (SAR) outside Earth orbit. The architecture of this radar instrument supports multiple polarimetric modes of operation (Putrevu et al. 2016, 2020) and illustrates the value of dual-frequency quad polarimetry for lunar and planetary applications. Regarding applications to the lunar surface, the main novelties provided by this instrument are the following:

**Frequency bands:** For the first time ever, it is possible to analyze L-band (~24 cm wavelength) polarimetric radar images of the Moon from the DFSAR instrument. Compared to previous ground-based (e.g., Arecibo, Goldstone) and orbital-based (e.g., Mini-SAR and Mini-RF) radar data of the Moon collected at the 70 cm (P band), 12.6 cm (S band), and 3.8 cm (X band) wavelengths, the new L-band SAR allows for a penetration depth of ~3 m (for dry, low-loss soils) with particular sensitivity to ilmenite content and surface or buried rocks with a diameter of about 1/10 the radar wavelength (i.e., ~2 cm) and larger. Additionally, combined with the S-band radar mode of operation (12 cm wavelength) along with data from the Mini-SAR and Mini-RF radars on board the Chandrayaan-1 and LRO missions respectively (Nozette et al. 2010), the DFSAR instrument can characterize the radar scattering properties of a top few meters of the lunar surface.

**Full-polarization:** The DFSAR architecture introduced the fully polarized (FP) case among polarimetric imaging radars for planetary missions, in which the intrinsic data product is the 4 × 4 scattering matrix of each resolved element in the scene. FP systems alternately transmit two orthogonal polarizations and record both received polarizations (HH, HV, VH, and VV) and allow much more information to be extracted from a scene compared with single- and dual-pol SAR data (Zyl et al. 1987). For each combination, the first letter refers to the transmitted polarization, while the second one refers to the received sense.
with H and V denoting horizontal and vertical, respectively. After applying certain symmetry relations, this can be reduced to a $3 \times 3$ array such as the compressed Sinclair matrix or the compressed Stokes matrix (Lee & Pottier 2009); such reduced forms are referred to as quadrature-polarimetric SAR (quad-pol for short). Polarimetric channels acquired in quad-pol mode maintain their relative phase, so they can be combined coherently to form new channels (Lee & Pottier 2009) or to compute statistical higher-order parameters by target decompositions (e.g., Cloude & Pottier 1996). Note that in the polarimetric radar literature, the terms “FP” and “quad-pol” are often used synonymously, so we adopt similar notation throughout this paper. Quad-pol data produce a unique scattering matrix using any combination of transmitted and received orthogonal polarizations (e.g., linear, elliptical, hybrid), which permit analysis of the surface scattering behavior in all possible configurations of the transmitted and received signal polarization. Such uniqueness does not apply to single, dual, and compact polarimetric SAR configurations (e.g., Lee & Pottier 2009).

**Spatial resolution and look geometry:** To address a variety of scientific objectives, DFSAR collects data of the lunar surface at a wide range of slant-resolution options spanning 2–75 m per pixel. The high spatial resolution of DFSAR enables mapping of lunar craters and other geological features, especially in the polar regions, with finer details. Compared to previous orbital and ground-based radars, high-spatial-resolution DFSAR data also provide a sufficient number of samples for the investigation of distributed scatterers that are common in lunar geologic settings. Radar backscatter will also vary depending on the incidence angle of the radar beam. Previous studies using FP SAR data of terrestrial settings indicated that surface scattering models could be developed using the relationship between surface roughness and dielectric parameters and the intensity, polarization, and angular dependence of the backscattered wave (e.g., Campbell et al. 1993; Campbell & Shepard 1996; Campbell 2002). In addition, theoretical models of radar scattering predict that at the near-nadir regime (i.e., from normal incidence to $\sim 20^\circ$), the reflection is dominated by locally smooth, radar-facing facets of the surface, and for incidence angles beyond $\sim 30^\circ$, diffuse scattering from the small-scale roughness predominates (Hagfors 1964; Farr 1993). The DFSAR instrument is designed to collect data at a wide range of angles of incidence, from 9.6° to 36.9° (Putrevu et al. 2016, 2020)). This capability enables the acquisition of dual-wavelength ($L$- and $S$-band) radar images at multiple viewing angles to address some important applications such as the estimation of wavelength-scale physical properties (e.g., slope, surface height, slope; Campbell et al. 1993; Campbell & Shepard 1996), and the reduction of the ambiguities related to the interpretation of high circular polarization ratios (CPRs) exhibited by certain lunar surface features (Campbell 2012; Spudis et al. 2013; Fa & Eke 2018; Fassett et al. 2018; Virkki & Bhiravarasu 2019).

The text is organized as follows. Section 2 describes the scientific background and objectives of DFSAR investigations. Section 3 is dedicated to the instrument overview and calibration of DFSAR data products. A brief description of DFSAR polarimetric parameters along with their physical interpretation is discussed in Section 4. In Section 5, we discuss the DFSAR data used in this preliminary analysis. Results are presented in Section 6, followed by conclusions and future work in Section 7.

**2. Background of Investigation**

**Polar volatiles:** The nature and distribution of volatiles (e.g., water ice) at the permanently shadowed regions (PSRs) of the Moon has been a subject of considerable controversy, due to contrasting interpretations of the polarimetric behavior of the radar backscatter from these regions at near-zero phase (bistatic) angles (e.g., Nozette et al. 1996; Campbell et al. 2006; Spudis et al. 2010, 2013; Fa & Cai 2013; Eke et al. 2014; Fa & Eke 2018; Fassett et al. 2018). The presence of ice in the regolith causes a distinct, but not unique, CPR signature due to forward scattering of radar signals by cracks and voids in the ice, which has an intrinsically low microwave loss. Due to the low-loss properties of ice, CPR values above 1 have been observed for the Galilean satellites (Campbell et al. 1978; Ostro et al. 1992; Black et al. 2001) and the Greenland ice sheet (Rignot 1995). Within some polar craters of Mercury, features exhibiting a very strong radar backscatter and enhanced CPR values along with a high degree of correlation between radar-bright features and regions of permanent shadow have been attributed to several radar wavelengths-thick slabs of water ice (Harmon et al. 1994, 2001, 2011; Chabot et al. 2012, 2013, 2018). One of the most plausible mechanisms that could produce these high backscatter and CPR values is the coherent backscatter model (Hapke 1990), which requires scattering centers (cracks or inhomogeneities) embedded in a low-loss matrix such as ice (Hapke & Blewett 1991; Mischchenko 1992). Unlike the Galilean satellites and Mercury poles, radar observations of lunar PSRs did not show radar-bright regions that indicate the presence of large expanses of water ice but indicated that the possible ice could be present as few wt% in the uppermost meter of regolith or mixed as patches within the polar regolith as “dirty ice” (Nozette et al. 1996; Thompson et al. 2011; Thomson et al. 2012). If, however, the lunar ice is present in the regolith in the form ice-filling pores, then radar scattering differences might be too small to detect (Thompson et al. 2011; Virkki & Bhiravarasu 2019). Another line of evidence from the interpretation of Mini-SAR and Mini-RF radar data suggested that both poles contain abundant water in the form of relatively “clean” ice, within the upper couple of meters of the lunar surface (Spudis et al. 2010, 2013). Other work has called this interpretation into question and has interpreted PSRs with enhanced CPR signatures as indicative of effects due to roughness, not water ice (Simpson & Tyler 1999; Campbell et al. 2006; Fa & Cai 2013; Eke et al. 2014; Fa & Eke 2018).

**Impact ejecta and melts:** Polarimetric radar is a powerful tool for studying lunar impact craters, exposing their “rough” ejecta deposits and associated impact melts (Campbell 2012; Neish et al. 2014). New multiwavelength radar observations would enable constraining block sizes and to distinguish the surface from buried rocks in proximal and distant ejecta deposits (Campbell 2002; Ghent et al. 2016). The impact melt deposits have very high CPRs compared to other features on the Moon, suggesting that their surfaces are some of the roughest material on the Moon at the centimeter to decimeter scale, even though they appear smooth at the meter scale (e.g., Carter et al. 2012; Neish et al. 2014, 2017).

**Radar-dark halos:** A certain class of impact craters on the Moon have radar-dark halos around them, comprised of fine-grained, rock-poor ejecta distal to the blocky proximal ejecta (Ghent et al. 2005). The composition of this halo material is suggested to be a mixture of crater ejecta and preexisting regolith (e.g., Ghent et al. 2016), and radar data at multiple...
wavelengths could provide significant new insight into the behavior of ejecta and their effects on the surrounding terrain.

### 2.1. DFSAR Objectives

After the end of the commissioning phase in September 2019, the science mission phase of Chandrayaan-2 has started, which will continue until any major contingencies hamper the spacecraft operation and/or science data reception. The DFSAR instrument is designed to address key science questions related to the aforementioned target properties during the Chandrayaan-2 science mission phase, by probing the lunar regolith to understand surface physical and dielectric properties and structure, and by exploring the PSRs at the lunar poles. Specifically, DFSAR data will be utilized in characterizing the physical properties of lunar mare, impact craters, and their associated ejecta and melt deposits, volcanic deposits including pyroclastic mantling material and domes, buried layering, and volatiles at lunar poles.

The DFSAR architecture is new for planetary missions. The quad-polarity design (Sections 3 and 4) provides data sufficient to measure the $2 \times 2$ complex scattering matrix of the backscattered field, which in turn leads to the extraction of all information available in the radar reflections. The fully polarimetric radar scattering theory contains all of the scattering information for any arbitrary polarization state (e.g., circular, linear, and hybrid modes), thus providing better tools than CPR alone to develop scattering models of the lunar terrain (Fa et al. 2011; Campbell 2012). Moreover, measuring the full scattering matrix allows a powerful observation space sensitive to shape, orientation, and dielectric properties of the scatterers within a resolution cell (Henderson & Lewis 1998).

### 3. DFSAR Instrument Overview, Operations, and Calibration

The DFSAR instrument is designed with a configuration that enables operations in dual-frequency and quad- and compact-polarity imaging modes to collect information about the scattering properties of the lunar surface at multiple look-angle geometries. For a complete review of technical specifications, system design, and configuration of the DFSAR instrument, the reader is referred to Putrevu et al. (2016, 2020). In this section, we briefly describe the system configuration, followed by the calibration of DFSAR data products.

#### 3.1. DFSAR Instrument Overview

DFSAR is designed as two independent radar systems in the $L$ band and $S$ band sharing a common microstrip planar antenna. Such a configuration enables standalone $L$/$S$-band imaging as well as synchronous $L$ and $S$ imaging; the latter is ensured by having synchronization and timing signals sourced from $L$-band systems to those of $S$ band. Linear frequency modulated (LFM; or more commonly referred to as chirp) signals of requisite bandwidth (BW; selectable from 75 to 2 MHz) are transmitted through two transmit chains feeding to the H and V ports of the antenna. Backscattered signals from the target are collected by the antenna and passed through H and V chains of the receiver, followed by digitization, data processing, and formatting in the data-acquisition system to meet the data-rate constraints of the Chandrayaan-2 mission. Transmitters of each band are configured in two chains for H and V, with gallium-nitride-based solid-state power amplifiers (SSPAs) in a high-efficiency design. Each chain of SSPA has a six-bit digital phase-shifter to enable requisite phase-setting for transmitting in either circular or linear polarization. Receivers are configured with low-noise and high-gain amplifiers to meet the high signal-to-noise ratio requirements of the system. A key feature of the DFSAR system is its capability of onboard range compression implemented in the data-acquisition system and provides data-rate reduction to the tune of 70%. This compression technique is a major factor in realizing the FP mode of operations, which is data-rate and volume intensive (about twice that of a hybrid-polarization implemented in Mini-RF/Mini-SAR instruments). Table 1 lists the major specifications of the DFSAR system.

Onboard calibration pulses flank the imaging pulses (data-windows) and include chirp replica and noise pulses as pre- and post-imaging sequences. The replica signal refers to the signal that is coupled (−20 dB) at both ends of the transmitter and the receiver. The DFSAR instrument is configured on the Chandrayaan-2 orbiter with the planar antenna mounted on the moon-viewing panel (positive yaw) of the spacecraft cuboid, with its electronics packages integrated behind this panel. During lunar imaging, the side-looking geometry of the radar is affected by roll, tilting the orbiter by the required look angle.

#### 3.2. DFSAR Preflight (Lab) Characterization

Lab characterization is essential for any instrument to meet the performance specifications, more so in the case of a polarimetric imaging system like the DFSAR. The instrument has undergone extensive in-lab characterization to ensure high-quality performance in terms of radiometry as well as polarimetry. Some of the salient tests conducted are as follows:
1. Radio frequency (RF) measurements of transmit peak power, receiver gain and noise figure at different manual-gain control (MGC) settings, front-end losses, receiver flatness, and imbalances in amplitude and phase are some of the key parameters for DFSAR characterization. These parameters have been characterized over the operating temperature range of 0°C–45°C, at the individual subsystem/component level. Important RF parameters were additionally characterized during thermo-vacuum tests of integrated payload.

2. Receive chains have been characterized to generate mapping of the digitized signal counts to input signal power, which is an important input for radiometric calibration. The entire chain from the input antenna port through front-end elements and receiver is characterized as an integrated unit to yield the above mapping. Similarly, transmit replica calibration path is also characterized, which serves as a reference to track any drifts in transmit and receive gain.

3. Far-field measurements of the antenna patterns for both bands and H and V polarizations have been done in a compact antenna test facility (CATF). Important parameters derived from the patterns include antenna gain, beam squint angles, and cross-polarization ratio, which play an important role in establishing the calibration equation and for other postprocessing requirements.

The above measurements meet the design specifications of the instrument and directly feed to the radiometric calibration of the SAR data, as described in Section 3.4.1. The received counts are first converted into expected backscattered signal power at the antenna-port input and used further to derive sigma naught ($\sigma^0$) by applying the above-mentioned lab measurements. $\sigma^0$ is the normalized measure of the radar return from a distributed target and is defined as per unit area on the ground.

1. The DFSAR instrument is designed to operate in any of these modes: single (HH/VV), dual (HH+HV/VV+VH), hybrid (LHCP/RHCP), or Full pol (FP). However, FP is the mode being used for systematic acquisitions of the lunar polar regions. At select locations of the Moon, imaging using hybrid polarimetry is also being performed, for which LHCP (left-hand circular polarity on transmission) mode is the default setting. For polarimetric characterization of the transmit chains, axial ratio measurements are crucial. In addition to characterizing the quality of circularly polarized transmission using axial ratios (Figure A1 in the Appendix), the measurements have yielded the best phase settings of the six-bit phase-shifters for the LHCP, right-hand circular polarity (RHCP), and FP modes (Figure 1).

2. Orientation of left/right circular transmit polarization is determined using helical antennas of the corresponding polarization. This has been used as a foolproof method to verify the handedness of the circular polarization.

3. Amplitude and phase imbalances between H and V receive chains (covering the entire path of the antenna, front end, and respective receiver) for polarimetric performance.

4. Using an external optical fiber delay-line, end-to-end characterization of the payload was carried out. The external setup consisted of a horn antenna followed by a circulator, delay-line, and appropriate attenuator pads; the transmit signal is collected by the horn antenna and passed through the circulator to the delay-line; the delayed signal is fed back to the circulator, to enable its transmission toward the DFSAR antenna. Using this, the process of replica acquisition and chirp compression using onboard range compression was validated.

5. The performance of the onboard range compression was validated using simulated distributed target data. The range-expanded data was played back from the Arbitrary Waveform Generator to the DFSAR data-acquisition system for range compression using chirp replica from pre-imaging calibration. The onboard range-compressed data were evaluated for impulse response parameters with respect to that of conventional on-ground processed output and found to be meeting the performance requirements.

6. Another key characterization activity is ensuring the phase synchronization between the $L$ and $S$ bands to be better than 6°. Chirp replicas for $L$ and $S$ bands compensate for any initial bias in phases, thereby ensuring a phase synchronization of better than 6°, as measured and verified.

7. DFSAR has an experimental radiometer mode for which the requisite measurements, data acquisition, and analysis will be described in a separate paper. However, for the sake of completion, an important characterization deserving mention here is that of radiometer mode full-chain “hot-load” calibration performed using microwave absorbers. This calibration complements the “cold-sky” measurement in operational onboard conditions toward retrieval of brightness temperatures.

### 3.3. DFSAR Operations, Data Acquisition, and Product Generation

Chandrayaan-2 was launched using ISRO’s Geosynchronous Satellite Launch Vehicle (GSLV)-Mark III on 2019 July 22 from Sriharikota, the spaceport of India. It was inserted into lunar orbit on 2019 August 20 and was brought to a near-circular orbit of 80 km × 120 km, with the periapsis and apoapsis close to the lunar south pole and north pole, respectively. Subsequent to the initial commissioning phase activities of DFSAR during 2019 September, the instrument has been tasked to systematically cover the lunar poles (85° and poleward) and is making full-polarimetric measurements in the designated time slots (referred to as dawn–dusk seasons). Since then, DFSAR has been operational and beaming high-quality data to cater to various scientific studies. After this initial campaign, DFSAR will observe the lunar polar regions (±70° to ±90°) at the $L$ band along with selected nonpolar targets at the $L$ and/or $S$ bands. Based on this coverage, selected regions (e.g., where anomalous CPR values exist) will be observed with $L$- and $S$-band simultaneous, hybrid- and full-polarization imaging. Over the course of the mission, we will plan to acquire global $L$-band FP data of the lunar surface to integrate with other global lunar data sets. Raw data from the DFSAR acquisitions are downlinked at different ground stations and processed to generate products at various levels, which are denoted below:

1. **Level-0A/0B**: represent range-uncompressed and compressed, respectively. These are raw-data products with
ancillary data for further processing. In this, different processes (like block adaptive quantization) applied to reduce the data rate are reversed to represent the original instrument data.

2. **Level-1A**: represents seleno-tagged single-look complex (SLC) data.

3. **Level-1B**: represents seleno-tagged ground-range product.

4. **Level-2**: represents seleno-referenced product.

A Range–Doppler algorithm (Cumming & Wong 2005) is used for processing DFSAR stripmap data, which efficiently processes the SAR data in the frequency domain. The time-domain correlation operation is replaced by the frequency domain multiplication operation. Each pixel of a given slant image is tagged to a particular seleno-location, i.e., latitude/longitude, value using the slant range grid generated from the orbit attitude information, a digital elevation model (DEM) available at ~118.5 m spacing from the LRO Laser Orbital Laser Altimeter (LOLA; Mazarico et al. 2011) and radar parameters.

### 3.3.1. Image Quality Analysis

Any feature or target of dimensions similar to or smaller than that of the SAR resolution cell and with sufficient backscatter energy will appear as a “point target” in a SAR image. Such point targets are suitable candidates to characterize the impulse response of the SAR system. DFSAR data obtained during the commissioning phase have been processed and evaluated for payload performance. In this process, range and azimuth impulse responses have been analyzed for opportunity point targets, because no ideal point targets are available on the lunar surface. System impulse response characteristics such as peak-to-sidelobe ratio (PSLR) and resolution have been derived from these opportunity targets, and some of the sample results are shown in Table 2 and Figure 2. After the radiometric calibration process (described in Section 3.4.1) of DFSAR, noise-equivalent sigma naught (NESZ) was derived (Table 3) from the noise floor of the system, measured during onboard calibration data acquisition. NESZ represents the lowest $\sigma^0$ that can be measured by the instrument. Radiometric resolution has been estimated over homogeneous regions located in the processed images. All the DFSAR performance metrics are found to match the prelaunch expected values.

### 3.4. Calibration of the Data Products

As the DFSAR is a polarimetric system with the capability of acquiring quad-pol and compact polarimetric (CP) data, radiometric and polarimetric calibration are critical for the derivation of lunar surface physical parameters.
3.4.1. Radiometric Calibration

Radiometric calibration provides the coefficients to map SAR data (pixel values) to their corresponding backscatter coefficients, i.e., the normalized radar cross section of targets. Due to the absence of any calibration targets on the Moon with the well-known radar cross section in the L and S bands, the calibration is performed based on the radar equation, using lab-characterized system parameters:

\[ P_r = k_v \times (\text{DN})^2 \times g_{oa} \times g_{ot} \times \frac{L_b}{\text{SF}}, \]

where \( \text{DN} \) represents the digital number corresponding to a pixel after processing; \( k_v \) is a factor that accounts for quantization to convert count to power at the input of analog-to-digital converter with 50 ohm impedance; \( g_{oa} \) and \( g_{ot} \) are the pulse repetition frequency (PRF)/Doppler BW and the range sampling frequency/range BW corresponding to the over-sampling factors for azimuth and range respectively; \( L_b \) is beam-shape loss; and \( \text{SF} \) is the scale factor applied during onboard processing, for a given polarization chain.

Using the above equation, processor-gain-related factors are first corrected to convert the digital counts into raw power, as would be present at the receiver output. The raw power as computed from Equation (1) is then inverted using Equation (2) to derive \( \sigma^0 \), with the inputs from lab-measured parameters:

\[ P_r = P_t G_t G_r \frac{\lambda^2 \sigma^0}{(4\pi)^3 R^2 L S} \times \frac{c t}{2 \sin \theta} \times \frac{G_{rx}}{\text{MGC}_{rx}}, \]

where \( P_r \) is the power received corresponding to a digital count of a pixel; \( P_t \) is transmitted power; \( G_t \) and \( G_r \) are the transmit and receive antenna gains respectively; \( G_{rx} \) is the receiver gain; \( R \) is the slant range of the pixel location; \( \lambda \) is the radar wavelength, \( \theta \) denotes the incidence angle, and \( c \) is the velocity of light; \( t \) is the pulse width or chirp duration; \( \text{MGC}_{rx} \) is the attenuation setting by the manual-gain controller of the receiver; \( L \) is the receiver chain front end loss; and \( L_s \) is the azimuth antenna length.

The processed images are corrected using the lab-measured antenna patterns for radiometric normalization, so that a uniform/single calibration equation can be used for the entire data set, irrespective of the pixel position. In the case of Level-1A (SLC) products, the relationship between the digital number (DN) and the backscattering coefficient \( \sigma^0 \) for each polarization channel can be written as:

\[ \sigma^0_{ij} = \frac{\text{DN}_{ij}}{10(\text{dB})}; \quad i = 1, 2 \ldots L; \quad j = 1, 2 \ldots M, \]

where \( \text{DN} = I + Q_i \), and I and Q are, respectively, the real part and the imaginary part of the complex pixel value; \( K \) is the absolute calibration constant and is constant over all SLC products; and \( L, M \) are the number of lines and columns in the product. The calibration constant \( (K) \) is provided with the associated xml label file for each data set and is different for products Level-1A and Level-1B and above. For Level-1B and above products, \( \sigma^0 \) in linear scale to dB can be directly derived as

\[ \sigma^0_{[\text{dB}]} = 20 \log_{10}(\text{DN}) - K. \]

Specimen results as shown in Figure 3 and Table 4 compare the radiometry of S-band DFSAR hybrid-polarimetric mode data with S-band Mini-RF data over a part of the Gauss crater region that has overlap from both instruments. The radiometric accuracies (with incidence-angle-independent gamma-naught, \( \gamma^0 \)) obtained over different sample regions are found to be better than 2 dB. In the absence of any L-band data and any full polarimetry data of the Moon, validation could only be done for the process adopted for radiometric calibration, using S-band hybrid-polarity data of Mini-RF as reference. Considering that there are no additional processes that may hinder such a generalization for either L-band or full polarimetry data, a radiometric accuracy of better than 2 dB can be safely considered as a worst-case value.

3.4.2. Polarimetric Calibration for Full Polarimetry

Polarimetric characterization of a region of interest is heavily dependent on the quality of polarimetric calibration that
establishes the interchannel (HH/VV/HV/VH) phase relationships. Calibration errors can affect polarimetry products (coherency and covariance matrices and the other derived daughter parameters) thereby leading to misinterpretation of scattering/physical processes under study. The “measured” scattering matrix (Section 4.1) and the distortion matrices corresponding to the transmit [T] and receive [R] paths of the radar signal. Considering the order of interactions between the signal and target, the “measured” scattering matrix may be written as

\[
\begin{bmatrix}
S_{HH} \\
S_{HV} \\
S_{VH} \\
S_{VV}
\end{bmatrix}^{\text{measured}} = K(\gamma)[R]
\begin{bmatrix}
S_{HH} \\
S_{HV} \\
S_{VH} \\
S_{VV}
\end{bmatrix}^{\text{actual}} [T],
\]

with

\[
[R] = \begin{bmatrix}
 f_i(\gamma) & \delta_1(\gamma) \\
\delta_2(\gamma) & 1
\end{bmatrix}
\]

\[
[T] = \begin{bmatrix}
 f_i(\gamma) & \delta_1(\gamma) \\
\delta_2(\gamma) & 1
\end{bmatrix},
\]

where \(f_i\) and \(f_r\) represent channel imbalances in HH polarization with reference to that of VV, for the transmit and receive paths respectively. They comprise differences in antenna patterns between the polarizations and also the gain and phase of their respective paths. \(\delta_1\) and \(\delta_2\) represent the cross-talk terms HV and VH, respectively, mainly arising from the antenna parameters.
cross-polarization components. As we are dealing with imaging in lunar conditions with negligible ionosphere, Faraday rotation of the signals is not expected and thereby not considered in the above equation. After the SAR image formation and radiometric calibration, polarimetric calibration is applied to the resulting data. The methodology adopted for this calibration closely follows that of Sun et al. (2018) due to the following major factors: (1) the methodology does not depend on any external calibration targets (like corner reflectors) and is derived from the SAR data itself, which is a necessity in the case of lunar imaging. (2) Cross-talk estimation and channel imbalances are estimated in an iterative fashion due to their mutual dependency. (3) The computations are based on a derived covariance matrix, which results in more limiting equations to derive the phase estimates.

DFSAR polarimetric calibration is based on the above methodology while focusing on dominant surface and volume backscatter regions for co- and cross-polarization signal responses. The results in Figure 4 show phase histograms for uncalibrated and calibrated L-band (Figures 4(a)–(b)) and S-band (Figures 4(c)–(d)) acquisitions, corresponding to co- and cross-polarization cases. Phase corrections of the order of $-50^\circ$ for the co-polarization phase and $-5^\circ$ for the cross-polarization phase are required for phase calibration of DFSAR L-Band acquisitions. Corresponding numbers for the calibration of S-band FP acquisitions are $-100^\circ$ and $+100^\circ$ for the co- and cross-polarization phase, respectively. For L-band data, the stability of the co- and cross-polarization phases observed over multiple acquisitions (Figures 4(e)–(f)) is typically within $\pm 10^\circ$, indicating that a uniform set of polarimetric calibration factors can be applied for the DFSAR data sets. However, monitoring of polarimetric calibration would be continued for future acquisitions, and any deviations will be taken care of by estimating new calibration coefficients and incorporating them during the subsequent processing of the DFSAR data.

For a holistic view of the quality of the polarimetric calibration, opportunity point targets have been analyzed for polarimetric signatures, both in terms of polar coordinates (ellipticity angle and orientation angle) and Cartesian coordinates representing Stokes parameters (Raney 2007; Lee & Pottier 2009) over a Poincaré sphere. The results (Figure 5) show an excellent match with those of an ideal trihedral corner reflector (Zyl et al. 1987) and give the required confidence on the quality of the polarimetric calibration.
3.4.3. Polarimetric Calibration for Compact Polarimetry

The calibration approach for CP data is a subset of that adopted for calibration of full-polarimetric data. The channel and phase imbalances were derived from L- and S-band data over some uniform regions on the lunar surface. The relative phase histograms computed from Stokes parameters (Raney 2007; Raney et al. 2011) were derived over several CP data acquisitions to establish the consistency of amplitude and phase-calibration parameters. The order of relative phase (δ) corrections is about 40° for L-band and 10° for S-band data to achieve the phase calibration. The results of applying this correction is shown in Figure 6 in the form of uncalibrated and calibrated δ and degree of polarization (m) histograms for a set of DFSAR L- and S- band CP data.

4. DFSAR Data Analysis and Interpretation

4.1. Quad-polarity Data Parameters

In FP mode, the DFSAR transmits two orthogonal polarizations (Horizontal “H” and Vertical “V”) on a pulse-to-pulse basis and receives the scattered waves in two orthogonal polarizations (in the same basis as used for transmission). Thus, for one snapshot of the received echoes, we can construct a 2 × 2 complex scattering matrix [S] that relates the two-dimensional transmitted and received electric field vectors [E] as (Lee & Pottier 2009; Cloude 2009)

\[
\begin{bmatrix}
E_H \\
E_V
\end{bmatrix}_{\text{receive}} = \frac{\exp(-ikr)}{r} \begin{bmatrix}
S_{HH} & S_{HV} \\
S_{VH} & S_{VV}
\end{bmatrix}_{\text{transmit}} \begin{bmatrix}
E_H \\
E_V
\end{bmatrix}_{\text{transmit}},
\]

where all the elements are complex valued, and the subscripts H and V indicate horizontal and vertical polarization, respectively. The factor \(\exp(-ikr)/r\), where \(k = 2\pi/\lambda\) is the wavenumber, expresses the attenuation for a spherical wave of a radius that equals the distance (r) between the scatterer and the radar. The elements of [S] describe the amplitude and phase changes imposed on the incident field by the scattering process, and in monostatic configurations like that of DFSAR, [S] becomes symmetric, i.e., \(S_{HV} = S_{VH}\) for all reciprocal scattering media (Lee & Pottier 2009; Campbell 2002). Once the full scattering matrix is measured, any arbitrary complex scattering amplitude can be reconstructed as a linear combination of the elements of the measured scattering matrix.

In order to describe the polarimetric scattering behavior of the lunar surface, which is a distributed target, a second-order statistical formalism is required. The most common formalism to fully characterize distributed scatterers is the 3 × 3 coherency [T] or covariance [C] matrix defined by the outer product of a three-dimensional scattering vector in the Pauli or lexicographic formulation (for a review, the reader is referred to Lee & Pottier 2009). Followed by this, multilook processing (averaging several independent estimates of reflectivity) on [T] and [C] matrices is required to assess the scattering mechanism of distributed targets. This averaging process (of “N” looks) also reduces the speckle (by a factor of \(N^{1/2}\)) and degrades the azimuth resolution (by a factor of N). As the SAR image is formed by coherently processing returns from successive pulses, this results in a pixel-to-pixel variation in intensity, and the variation manifests itself as a granular pattern called speckle (Lee & Pottier 2009). A common approach to speckle reduction is to average several independent estimates of reflectivity, as described above. After calculating the multilooked [T] and [C] matrices and by assuming that there is no correlation between the cross-polarized (HV, VH) and like-polarized (HH, VV) linear components for natural surfaces (Campbell 2012), it is straightforward to calculate the backscatter coefficients (\(\sigma^b\)) in the two
circularly polarized components from the linearly polarized terms as given below (Campbell 2002, 2012):

\[
\text{Same sense Circular } (\sigma^\text{SC}) = \frac{1}{4} \left[ \sigma^\text{HH} + \sigma^\text{VV} + 4\sigma^\text{HV} - 2\text{Re}(S_{\text{HH}}S_{\text{VV}}^*) \right]
\]

\[
\text{Opposite sense Circular } (\sigma^\text{OC}) = \frac{1}{4} \left[ \sigma^\text{HH} + \sigma^\text{VV} + 2\text{Re}(S_{\text{HH}}S_{\text{VV}}^*) \right]
\]

The CPR, which provides significant information on the physical characteristics of the target surface (e.g., Campbell et al. 1993; Harmon et al. 1994; Rignot 1995; Harmon et al. 2001; Campbell 2012), can thus be calculated using

\[
\text{CPR} = \frac{\sigma^\text{SC}}{\sigma^\text{OC}}.
\]

Scattering decompositions are widely applied over polarimetric radar data for interpretation, physical information extraction, segmentation, and/or as a preprocessing step for geophysical parameter inversion (e.g., Cloude & Pottier 1996, 1997; Lee & Pottier 2009; Cloude et al. 2012; Raney et al. 2012). In general, the decompositions of second-order scattering matrices (i.e., [T] or [C]) are rendered into two classes: eigenvector- and eigenvalue-based decompositions and model-based decompositions (Cloude & Pottier 1996). For the lunar surface, it is of interest to generate the concept of an average or dominant scattering mechanism for the purposes of classification or inversion of scattering data through such target decomposition theorems. For DFSAR quad-pol data results shown in Section 6.1 of this paper, we adopted two widely used decomposition methods, which are the eigenvector- and eigenvalue-based (Cloude & Pottier 1996) and Yamaguchi 4 component decomposition (Yamaguchi et al. 2005) methods. Specifically, we derive two important statistical parameters arising directly from the eigenvalues of the coherency matrix, which are the polarimetric scattering entropy ($H$) and the alpha angle ($\alpha$) to extract physical information from the observed scattering of microwaves by the lunar surface. They are calculated for the backscattering case as follows (Cloude & Pottier 1996):

\[
P_i = \frac{\lambda_i}{\sum_{j=1}^{3} \lambda_j},
\]

\[
H = -\sum_{i=1}^{3} P_i \log_2 P_i,
\]

where $0 \leq \lambda_3 \leq \lambda_2 \leq \lambda_1$ are the real nonnegative eigenvalues of $[T]$ and $P_i$ expresses the appearance probability for each contribution. For each eigenvector, the roll-invariant scattering alpha angle is calculated as

\[
\alpha = \sum_{i=1}^{3} P_i \alpha_i.
\]

The ranges of entropy ($0 < H < 1$) can be interpreted as a measure of the randomness of the scattering process. An entropy of $H = 0$...
indicates a non-depolarizing scattering process described by a single scattering matrix, while the other extreme, an entropy of \( H = 1 \), indicates the presence of a random noise scattering process, which depolarizes the incident wave completely regardless of its polarization (Lee & Pottier 2009). The average alpha angle ranges between 0° and 90° and is associated to the type of corresponding scattering mechanism. In general, \( 0° \leq \alpha \leq 30° \) corresponds to surface scattering processes, \( 40° \leq \alpha \leq 50° \) to dipole-like scattering behavior, and finally \( 60° \leq \alpha \leq 90° \) indicates dihedral or helix type scattering mechanisms (Lee & Pottier 2009).

### 4.2. Compact-polarity Data Parameters

While the FP SAR data mode in DFSAR provides optimum performance in target characterization due to its complete radar target information content, compact polarimetric (CP) SAR mode offers more information than acquired from single or dual-polarized SAR mode (e.g., Raney et al. 2011). The hybrid polarimetry mode of DFSAR follows from the Mini-SAR and Mini-RF radar configuration (Nozette et al. 2010), which receives orthogonal linear polarizations while transmitting circular polarization. The hybrid-polarimetric architecture of the DFSAR is a form of compact polarimetry (Raney et al. 2011), and the resulting data from this mode support calculation of the 2 × 2 covariance matrix of the backscattered field, from which follow the four Stokes parameters required to calculate \( \sigma_{oc}^{2} \), \( \sigma_{sc}^{2} \), and CPR. For brevity, we do not elaborate here the derivations of these parameters including the m-chi decomposition technique that is used for the analysis of DFSAR CP data (for a review, the reader is referred to Raney et al. 2012).

### 4.3. Image Artifacts

Some DFSAR images contain vertical streaks near the crater rims, with a spread in the azimuth direction. This has been observed most prominently at the far wall (from the radar’s perspective) of large impact craters that perfectly align (i.e., oriented perpendicularly) with the incident radar signals, giving rise to saturation of the signal, which after SAR processing appears as streaks (e.g., see the red vertical streaks in Figure 8). This is similar to the “cardinal effect” observed commonly from urban regions in terrestrial radar images due to the tendency of a radar to produce very strong echoes from a city street pattern or other linear feature oriented perpendicular to the radar beam (e.g., Henderson & Lewis 1998).

Other artifacts in the form of irregular patches observed in some DFSAR FP images are a manifestation of azimuth ambiguities (e.g., see the red irregular patches in Figures 8 and 11(a)). These are a result of the target returns that enter the system through azimuth side lobes of the antenna pattern and are aliased onto the processing (Doppler) BW. Though the ambiguity levels are quite low and visible in only FP co-pol backscatter images (HH, VV) after stretching/scaling, these artifacts are emphasized greatly in the decomposition images. One of the ways to reduce this azimuth ambiguity is to further increase the PRF, which is already doubled due to FP SAR requirements. However, increasing the PRF would have challenges in accommodating the data window and can lead to range ambiguities. Also, this would be only possible in new data acquisitions with new PRF settings. Hence, we have very recently adopted an alternative method, which is to process the (already acquired) data with reduced processing BW (say, half) and thereby excluding the aliased component of the returns. The reduced processing BW impacts the single-look azimuth resolution; however, in the case of the DFSAR, it can be taken care of by adjusting the corresponding multilook factor. We have tested this processing method on one of the FP data sets, and significant improvements are observed in the sense that most of these azimuth-ambiguity-related artifacts are removed. After testing this method over several data sets, we plan to process all the existing DFSAR FP data through this new processing chain.

### 5. Data Used and Study Area

As the DFSAR system probes the lunar regolith at two frequencies (L and S bands), it provides additional information on the physical properties of the upper meter or two of the lunar surface, as we show with some examples in the following section. Here, we present the first results of a campaign designed to evaluate the potential of quad- and compact polarimetry at the L and S bands using the new DFSAR data. To this end, several images over polar regions (~40% of 80°–90° latitude) and a few nonpolar regions have been acquired through 2020 October, providing data at different sensor configurations (polarization channels and incidence angles). Figure 7 shows DFSAR data coverage in both polar regions of the Moon. For this preliminary analysis, to characterize the crater ejecta and floor materials at the L- and S-band wavelengths, we sampled three different regions from the polar and nonpolar regions of the lunar surface. We restricted our analysis to crater interiors and their continuous ejecta blankets to isolate the response of the ejecta from surrounding materials. Typically, this (ejecta) region encompasses surface material within approximately 1–2 crater radii of the crater rim (e.g., Thompson et al. 1981).

We analyzed L-band FP mode data of secondary craters on the floor of the Peary crater (88.6°N, 33°E) in the north polar region, L- and S-band (dual-frequency) CP data of the Byrgius C crater (21.17°S, 64.52°W) located between Mare Humorum to the west and the Orientale impact basin to the east, and L-band CP mode data over a region near Manzinus C crater (70°S, 21.6°E) in the south polar region. Because each of the above data sets are acquired at different configurations, we processed them with different numbers of looks and at different resolutions, which are explained in the following section. For each case (polarimetric mode of acquisition), we focus on comparisons between DFSAR data and S-band zoom mode (Raney et al. 2011) data obtained from the Mini-RF instrument on board LRO. For context and comparison, we utilized optical images obtained from the Lunar Reconnaissance Orbiter Camera (LROC) Wide Angle (WAC) and Narrow Angle Camera (NAC; Robinson et al. 2010).

### 6. Results

#### 6.1. L-band FP Data

L-band data of the Peary crater region was acquired in FP mode at an incidence angle of 26° with a resolution (pixel spacing) of ~0.6 m in azimuth and 9.6 m in range. From the calibrated SLC data, we generated the 2 × 2 scattering matrix followed by the 3 × 3 [C] matrix (Section 4.1). The averaged [C] matrix was then used with the slant range grid file (Section 3.3) to project the data onto a lunar coordinate system.
Figure 7. Data acquisition by Chandrayaan-2 DFSAR over the (a) north and (b) south pole regions of the Moon. Both L- and S-band DFSAR image swaths acquired in different polarimetric modes (as indicated in the legend below) are overlaid on LROC WAC polar mosaics for context.

Figure 8. Peary crater (78 km diameter; 88.6°N, 33°E) in the lunar north pole region as viewed by L-band DFSAR in Yamaguchi 4 component decomposition images. The terrain-corrected DFSAR decomposition images are mosaicked and overlain on an LROC WAC North Pole mosaic. The white dashed polygon outlines the approximate extent of the Peary crater rim, and the polygons in magenta are the regions of permanent shadow as obtained from LRO’s LOLA instrument. The two “anomalous” craters and a “normal” fresh crater (marked with numbers “1”, “2,” and “fresh crater,” respectively) analyzed in this study are shown with arrows in yellow and are also separately shown in Figures 10 and 11. The color wheel highlights the colors for each scattering regime (red: even bounce; blue: single (odd) bounce; green: volume scattering). Note that the vertical lines and small patches in red near crater rims and edges of DFSAR image strips that appear here are artifacts as described in Section 4 and are not to be confused with the dominant scattering mechanism from those regions.
with a pixel spacing of 25 m pixel$^{-1}$. This yielded a $\sim$38 look average for each sampled location in an observation and an approximate $1/N^{1/2}$ ($N =$ number of looks) uncertainty in the CPR measurements of $\pm$0.16 (Campbell 2002). The process of transforming the unprojected data to a lunar grid also involved correcting the data for terrain height variations that are inherent in the oblique viewing geometry of the SAR instrument. The terrain correction of DFSAR data is performed by utilizing Chandrayaan-2 spacecraft orbit information and (per pixel) topography information available from a global digital elevation model (DEM) at $\sim$118 m spacing available from the LOLA instrument (Mazarico et al. 2011) on board the LRO mission. Finally, the multilooked, terrain-corrected [C] matrix is used to generate the backscatter coefficients, CPR, entropy, average alpha angle, and decomposition images (Figure 8) using Equations (7)–(12), required for our analysis. We analyzed two impact craters on the floor of the Peary crater that were classified as "anomalous" in previous studies (Spudis et al. 2010, 2013; Fa & Eke 2018; Virkki & Bhiravarasu 2019) along with a "normal" fresh, young crater to characterize their scattering properties at $L$ band. Anomalous craters are those having a high CPR in their interior but not exterior to their rims (Spudis et al. 2010, 2013; Fa & Eke 2018). All three impact craters are within PSRs at the north pole. Figure 9 shows histograms of CPR, wave entropy, and mean alpha angle plotted from the polygons interior and exterior ($\sim$1 crater radii) to the crater rim of these three craters as shown in Figure 11(b).

From Figure 9 (and Table A1 in the Appendix), we observe that the two anomalous craters (earlier observed at the $S$ band using Mini-RF data) appear anomalous at the $L$ band also, and the mean CPR and alpha angle values for the young, fresh crater region are relatively higher than those observed for both anomalous craters. The ejecta of young, fresh craters often stand out among other lunar features when observed at radar wavelengths (e.g., Campbell et al. 2010; Campbell 2012; Raney et al. 2012; Saran et al. 2012). This is due to the presence of a relatively high population of wavelength-scale scatterers found within the ejecta deposits. This roughness signature appears radar bright in the backscattered power images, associated with high CPR values that exceed unity (Campbell 2012). Also, rough surfaces are dominated by multiple-bounce backscattering, which randomizes the polarization (Campbell 2002, 2012). For this reason, the proximal ejecta region of the fresh crater appears distinct in the HV polarization (depolarized) image compared to HH and VV polarization images (polarized), shown in Figure 10. Apart from their difference in absolute values (shown in dB scale in Figure 10), the behavior of anomalous crater regions in all the three polarization channels (HH, HV, VV) appears to be similar.

In contrast with the anomalous craters, the distribution of mean CPR, entropy, and alpha angle at the interior and exterior of the fresh crater are almost identical, indicating that the roughness of the two areas is, on average, the same on the scale of centimeters to decimeters. Higher values of mean entropy ($\sim$0.87) and alpha angle ($\sim$41$^\circ$) clearly indicate the presence of a dominant volume scattering component at regions both inside and outside of the fresh crater. This can be observed qualitatively from Figure 11,
which shows the relative contributions of dominant scattering mechanisms for each crater region. In addition, we observe a nonzero negligible contribution from the double-bounce scattering component (~4 dB less than the lowest of the other two scattering components) and almost no contribution from the helix component of the Yamaguchi 4 component decomposed image of these crater regions. So we ignored the contribution of these two components in this analysis. An interesting observation from the relative contribution of the mean surface scattering component from Figure 11(a) is that the surface scattering distributions of scatterers for all the three crater regions appear similar. As a first approximation, this behavior may indicate that the depolarization of linear-polarized illuminating signals occurs likely through single scattering by rock edges or other discontinuities, rather than as solely multiple-scattering effects predicted by some analytical models. However, additional terrain parameters such as the rms slope (horizontal scale dependent) and dielectric properties are required to quantitatively validate this interpretation.

We then compared our observations with those of S-band Mini-RF data of these crater regions in Figures 11(c)–(d). The Mini-RF level-2 data of this region has a fixed look angle of 49° and is reduced to a pixel spacing of 15 m before analysis, giving a speckle noise of $1/N^{1/2} \approx 19\%$ for each pixel. We determined the mean CPR for interior and exterior regions of the three crater regions (from the polygons shown in Figure 11(d) and values given in Table A1 in the Appendix) and plotted them against those from DFSAR data for comparison, as shown in Figure 12. This scatter plot reveals an interesting observation—while the mean CPR values for fresh crater interior and ejecta are identical at the $L$ and $S$ bands, they are different for the anomalous craters analyzed in this study. For the anomalous crater regions, the $S$-band CPR is higher than that at the $L$ band, implying differences in the material properties and/or rock sizes present in those regions. The difference in the incidence angle between Mini-RF (49°) and DFSAR data (26°) does not fully explain this observation. Variations in CPR values for lunar craters are attributed to a combination of radar incidence angle, age, and target materials. For a given crater region, the other two parameters are the same, when comparing the CPR response solely as a function of incidence angle. Lunar surface slopes, such as those resulting from crater walls, can change the received radar echoes and CPR values because the surface slope affects both the local incidence angle and polarization state of the incident radar wave in the local coordinate frame (e.g., Fa et al. 2011). So, larger variations in CPR (as a function of incidence angle) are expected within crater interiors, owing to greater variations in the local slopes, compared to their exterior regions. On the contrary, we observed from Figure 12 that CPR differences for the exterior regions of the two anomalous craters are relatively higher compared to those of their interiors when observed at the incidence angles of DFSAR (26°) and Mini-RF (49°). Therefore, we attribute this result to the change in size distributions of scatterers as a function of wavelength ($L$ versus $S$ band) of observation. If rock size is the dominant effect causing such a discrepancy, radar signals at different wavelengths could be utilized to constrain their sizes and to distinguish the surface from buried rocks in the proximal ejecta of craters (e.g., Ghent et al. 2016). For the fresh crater region, similar CPR values at both $L$- and $S$-band indicate the presence of volumetric scatterers or roughness at centimeter and larger scales. In the case of the anomalous crater regions, the rocks responsible for $S$-band enhancement do not cause a similar enhancement in the $L$-band data, indicating that they are too small to be resolved by the longer-wavelength $L$-band signals. These observations imply that the material evolution of the interior and ejecta of anomalous craters are decoupled, supporting previous interpretations of a large set of lunar craters using $S$-band Mini-RF data (Fa & Eke 2018; Fassett et al. 2018). Interpreting this decoupling phenomena and CPR enhancement seen within the anomalous crater interiors is not the major focus of this work and will be addressed in a future paper.

6.2. Dual-frequency ($L$- and $S$-band) CP Data

We obtained simultaneous dual-frequency ($L$- and $S$-band) CP radar data of the crater Byrgius C, which was also classified as “anomalous” in previous studies (Fa & Cai 2013; Spudis et al. 2013). These data were acquired at an incidence angle of 24° and has a pixel spacing of ~0.5 m in azimuth and 9.6 m in range. From the calibrated SLC data, we generated the 2 × 2 covariance matrix for each frequency band and performed multilooking (in the azimuth direction) with $N = 46$ to obtain a slant range pixel spacing of 10 m in range and 24 m in the azimuth direction. We then used this multilooked 2 × 2 $[C]$ matrix to generate Stokes parameters followed by the CPR and $m$-chi decomposition images for $L$- and $S$-band data (Raney et al. 2012). For these data, the approximate uncertainty in the CPR measurement is ±0.15 for each sampled location.
Similar to the analysis presented in Section 6.1, we compared our DFSAR observations with those obtained from S-band Mini-RF level-2 data of the Byrgius C crater region (Figure 13). We observe that the mean interior CPRs for Byrgius C exceed 1 and are independent of wavelength (L vs. S band) and incidence angle (24° for DFSAR vs. ~50° for Mini-RF), contrary to the CPR trends observed from anomalous craters on the floor of Peary crater. The histograms of mean CPRs (Figure 14) also indicate that the average behavior over the crater interior is CPR ∼1 and is limited to less than ∼2.5. The mean CPR values (followed by the median value in square brackets) determined for the crater interior from DFSAR are 1.17 [1.14] and 1.15 [1.11] for the L and S bands respectively, compared to 1.16 [1.02] obtained from Mini-RF at S band. For the ejecta regions, they are 0.45 [0.42] and 0.7 [0.61] from DFSAR (at both L and S bands) and Mini-RF, respectively. The relatively higher differences between the mean and median values of CPR obtained from Mini-RF data of this region result in a more skewed distribution, as seen in Figure 14(c). Even after spatially averaging over a large number of looks to limit speckle effects, there are significant areas of the crater interior with CPR ratios greater than unity observed from both DFSAR and Mini-RF data. These characteristics are consistent with scattering by a combination of the quasi-specular (or small facet) and diffuse mechanisms (Campbell et al. 1993; Byrgius C 2002), most possibly due to the presence of a large number of wavelength-sized blocks and/or cracks present at the Byrgius C crater interior. The highest CPR enhancements (>2) may be attributable to a dihedral mechanism or double-bounce scattering from the randomly oriented dipoles (Campbell 2012). The m-chi decomposition images of the crater interior (Figures 13(b), (d), and (f)) indicate a mixture of volume and double-bounce scattering at work and support this interpretation.

6.3. L-band CP Data

For this analysis, we obtained very high-resolution (pixel spacing of ~0.4 m in azimuth and 1.8 m in range) L-band CP data of the Manzinus C crater region at an incidence angle of 30°. From the calibrated SLC data, we generated the 2 × 2 covariance matrix and resampled them using the slant range grid file to project the data onto a lunar coordinate system with a pixel spacing of 15 m pixel⁻¹. Unlike the previous two comparisons shown in Sections 6.1 and 6.2 in which we used level-2 Mini-RF data, we examined Mini-RF data of this region by processing the level-1 data using USGS ISIS3 software as described in Fa & Eke (2018) and Virkki & Bhiravarasu (2019). From the original spatial resolution of 7.5 m pixel⁻¹ of a level-1 product, we generated the orthorectified Stokes vector at 15 m pixel⁻¹ resolution to match that of the DFSAR L-band data. Finally, the terrain-corrected Stokes vector from both data sets is used to calculate the CPR and m-chi decomposition images as described in the previous section. Due to the fewer number of looks, for each pixel in both DFSAR and Mini-RF data used here, the uncertainty in the CPR measurements is on the order of ±0.2.
Figure 15 shows the $m$-chi decomposition images obtained from $L$-band DFSAR and $S$-band Mini-RF data, compared with a corresponding LROC NAC image. We selected three similar-sized ($\sim 0.8$ km) craters and marked them as “fresh,” “partially degraded,” and “highly degraded,” based on their roughness properties, evident from a visual inspection of the radar data (Figure 15). We determined the mean CPR for interior and exterior regions of the three craters from the $S$-band Mini-RF and $L$-band DFSAR and plotted them as shown in Figure 16. From Figures 15 and 16, we observe that the fresh crater appears distinct at both wavelengths from the other two degraded craters, as the radar signals are sensitive to scattering from the surface and buried rocks (size ranges of about 1/10 to 10 wavelengths) and roughness.

A radar-bright ejecta blanket that is excavated by the impact, which is not visible in the corresponding LROC NAC image (Figure 15), surrounds the fresh crater. For the partially degraded crater, which has CPR values intermediate to those of the fresh and highly degraded craters, this radar-bright ejecta halo is not distinct. The CPR values of both the crater interior and proximal ejecta of the highly degraded crater are comparable to the background regolith. CPR measurements obtained at multiple radar wavelengths often provide information about how lunar craters evolve with time, and how rocks at the lunar surface and in the near subsurface break down (e.g., Ghent et al. 2016; Fassett et al. 2018). A recent model using the $S$-band CPR and depth/diameter ratio of a set of simple craters suggests that the crater degradation state can be used to estimate their age (Fassett & Thomson 2014). One of their results suggested that the extent and magnitude of CPR enhancement decline as craters age. For the three crater regions, it is interesting to note that the mean CPR values at their corresponding interior and exterior regions at the $L$ and $S$ bands are almost similar (Figure 16), given the differences in wavelength and incidence angle. Unlike the results obtained from the comparison of the north polar craters as discussed in Section 6.1, the similar CPR values obtained from these three different crater regions suggest the presence of centimeter-to-decimeter blocks. However, the role of the polarimetric basis of observation (FP versus CP) in determining the CPR needs to be
further investigated to fully understand these observed CPR patterns. Nevertheless, if their ages can be calculated by other means (e.g., modeling and crater counting statistics), the distinct CPR patterns between the interior and exterior of these three craters may provide insights into how the rockiness of craters evolves with time.

Figure 13. Byrgius C crater region observed from the DFSAR and Mini-RF radar instruments. (a) CPR and (b) m-chi decomposed image derived from S-band Mini-RF CP data. (c) CPR and (d) m-chi decomposed image derived from S-band DFSAR CP data. (e) CPR and (f) m-chi decomposed image derived from L-band DFSAR CP data. In the m-chi images, red, green, and blue colors represent the double-bounce, random-polarized (volume scattering), and single-bounce component of backscatter respectively. North is up, and the radar look direction is indicated in each set of images. All the CPR images are stretched to the same color scale and overlaid on the total backscattered power images. The polygons shown in (a), (c), and (e) are used to generate the histograms shown in Figure 14. Also note the characteristic red halo (Raney et al. 2012) observed at the far wall (from the radar’s perspective) of the crater at the bottom of the Mini-RF m-chi decomposed image (b) that is missing in the corresponding DFSAR images ((d) and (f)) due to the change in the incidence angle of observation (i.e., 49° of Mini-RF vs. 26° of the DFSAR).

Figure 14. CPR histograms for interior and exterior regions of Byrgius C crater derived from (a) S-band DFSAR CP, (b) L-band DFSAR CP, and (c) S-band Mini-RF data.
7. Conclusions and Future Scope

DFSAR offers global radar data of the Moon at a range of spatial resolutions to investigate PSRs at the poles and to probe the lunar regolith to reveal their surface and near-subsurface physical properties by integrating with other global lunar data sets. The data calibration analysis presented here provides a means to compare on-orbit performance with prelaunch measurements, and the results are found to be matched with the prelaunch expected values, meeting their design requirements. By exploiting the full range of information contained in the FP SAR format, our results obtained in this preliminary analysis to decompose lunar radar echoes into physically meaningful components using their polarization properties help to understand the types of surfaces that produce different combinations of polarimetric behaviors. The radar polarization properties of a set of impact craters and their proximal ejecta at different wavelengths provide important information about their physical properties. We observe that craters in both PSRs (Peary secondary craters) and non-PSRs (Byrgius C) that are classified as CPR anomalous in previous S-band radar analyses appear anomalous at L-band also. In the case of the Byrgius C crater, we compared our dual-frequency results with the predicted behaviors of scattering by a combination of the small facet and diffuse mechanisms (Campbell et al. 1993; Campbell 2002), and the highest CPR enhancements (>2) are attributed to the dihedral mechanism (or double-bounce scattering) from randomly oriented dipoles (Campbell 2012). Based on our L- and S-band data analysis of crater freshness, further investigation of differences in the CPR evolution of different crater sizes would be a fruitful avenue for future work.

Analysis of FP radar data obtained from DFSAR along with other available radar data sets from the Mini-RF radars (S and X bands) and Arecibo Observatory (S and P bands) would enable us to better understand the variability in radar echoes from lunar regolith as a function of wavelength. A major advantage in utilizing the longer-wavelength L-band radar is its ability to penetrate the shallow (few meters at most) regolith, based on the loss tangent, to estimate the influence of regolith composition and (under favorable conditions) the nature of the substrate on the radar echoes. Moreover, the lunar water ice controversy cannot be resolved with existing monostatic data (prior to DFSAR), and new radar measurements are required. The unique dual-wavelength fully polarimetric DFSAR data are better poised to reduce the ambiguities associated with radar detection of water ice at the lunar PSRs by detecting distinct polarimetric signatures of “dirty ice” regions. Toward this end, constraining the distribution of dielectric properties and surface roughness across the lunar surface, and hence their role in radar surface scattering and subsurface penetration depth, is essential to reducing uncertainties associated with the identification and quantification of potential ice inclusions in the regolith. As linear polarization can provide a quantitative approach to understanding the surface physical properties (i.e., dielectric and textural properties; e.g., Campbell et al. 1993; Campbell 2002), we plan to derive the regolith dielectric constant from DFSAR FP data, followed by utilizing lab-based and theoretical dielectric mixing models (with several ice–regolith mixture combinations) to detect the possible presence of ice inclusions. Furthermore, surface roughness at the PSRs could be estimated by utilizing Mini-SAR, Mini-RF, and DFSAR radar data in conjunction, because all these data sets are acquired at different wavelengths and viewing geometry, and therefore offer further validation of centimeter to decimeter near-surface roughness.

Finally, with fully polarimetric SAR data of the lunar surface now available for the first time ever from the DFSAR, another significant area of research that could be explored is to investigate whether there is a fundamental difference in the capabilities of different radar architectures (e.g., FP versus CP) when it comes to measuring the polarimetric properties of a surface, as explored previously in Carter et al. (2017). Similar to the analysis presented in Carter et al. (2017), we plan to attempt the following comparisons of monostatic radar data by acquiring new DFSAR data sets (at similar viewing geometry).
in the near future: (i) DFSAR FP and CP data obtained at the same wavelength, (ii) DFSAR S-band FP and CP data with those of the Mini-RF S band, and (iii) DFSAR S-band FP and CP data with those of the Arecibo S band.

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Appendix
Supporting Information

This appendix contains a figure (Figure A1) and a table (Table A1). Figure A1 shows axial ratio measurements of the DFSAR instrument in the anechoic chamber as a part of its preflight characterization. Table A1 shows mean values of polarimetric parameters derived from DFSAR L-band fully polarimetric data of secondary craters on the floor of Peary crater near the lunar north pole.
Figure A1. Axial ratio measurements in the anechoic chamber in the fully integrated condition of the payload over the spacecraft (shown on the right of the figure). Toward the left is a linearly polarized horn antenna used to measure the signal intensity with different orientation angles.

Table A1
CPR, Entropy, and Mean Alpha Angle Values Derived from DFSAR L-band FP Data of Secondary Craters on the Floor of the Peary Crater Region

| Region          | CPR (DFSAR) | CPR (Mini-RF) | Entropy | Mean Alpha (degree) |
|-----------------|-------------|---------------|---------|---------------------|
| Fresh Crater    |             |               |         |                     |
| exterior        | 0.72 ± 0.14 | 0.68 ± 0.10   | 0.86    | 38.8                |
| interior        | 0.82 ± 0.16 | 0.87 ± 0.14   | 0.89    | 42.5                |
| Anomalous       |             |               |         |                     |
| exterior        | 0.30 ± 0.08 | 0.49 ± 0.09   | 0.62    | 23.1                |
| interior        | 0.70 ± 0.13 | 0.71 ± 0.15   | 0.85    | 38.7                |
| crater 1        |             |               |         |                     |
| exterior        | 0.33 ± 0.09 | 0.48 ± 0.09   | 0.63    | 24                  |
| interior        | 0.64 ± 0.13 | 0.67 ± 0.15   | 0.83    | 36.9                |

Note. For comparison, mean CPR values derived from corresponding LRO Mini-RF S-band data are listed here. These values represent the average of profiles analyzed from the regions interior and exterior (~1.5 radii) to the crater rims indicated in Figure 12, with 1σ standard deviation given as error for CPR.

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