Quantum Measurements, Nonlocality and the Arrow of Time

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A theory is developed which attempts to reconcile the measurements of nonlocal quantum observables with special relativity and quantum mechanics. The collapse of a wave function, which coincides with a nonlocal measurement by some macroscopic measuring device, is associated with the triggering of an absorber mechanism due to the interaction of the apparatus with the charges in the rest of the universe. The standard retarded electromagnetic field plus radiation damping is converted, for a short time during the collapse of the wave function, to an advanced field plus radiation. The reversal of the arrow of time during the wave function reduction permits communication in nonlocal quantum experiments at the speed of light, resolving paradoxes associated with measurements of correlated quantum states and special relativity. The absorber mechanism and the advanced field solution are consistent with conventional Friedmann-Robertson-Walker expanding universes.

I. INTRODUCTION

A recent Franson-type test of the Bell inequality [1,2] has demonstrated quantum correlations over more than 10 kilometers for energy and time entangled photons using a telecom fiber network [3,4]. This experiment reveals the nonlocal nature of quantum mechanics (QM) over large distances. It is predicted by quantum mechanics that this nonlocality will hold for the entire universe. The number of experiments of this kind, including experiments on photon tunneling [5], appears to show that nonlocal correlations are a permanent feature of QM.

The counter-intuitive predictions of QM have puzzled physicists since the formulation of the theory. The peculiar nonlocal effects arise from particles in entangled states as predicted by the linear superposition principle of QM:

\[ |\psi\rangle = \frac{1}{\sqrt{2}} (|\alpha\rangle_1 |\beta\rangle_2 + |\gamma\rangle_1 |\delta\rangle_2), \]

where \(|\alpha\rangle_1\) and \(|\gamma\rangle_1\) (\(|\beta\rangle_2\) and \(|\delta\rangle_2\)) are orthonormal vectors in Hilbert space for particle 1 (particle 2).

Although QM agrees remarkably well over a wide range of experiments, there has been an on-going consistency problem with relativity theory. The standard point of view of “pragmatic” practitioners of QM is that there is no violation of causality in nonlocal QM measurements. It is stated that quantum mechanics does this nonlocality will hold for the entire universe. The number of experiments of this kind, including experiments on photon tunneling [5], appears to show that nonlocal correlations are a permanent feature of QM.

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The “orthodox” view is that the quantum state vector (wave function \(\psi\)) characterizes the individual system completely, a point of view tenaciously opposed by Einstein [3] and championed by Bohr [4]. Bohr argued that it is impossible to make a sharp separation between the behavior of atomic objects and the interaction with the measuring apparatus which defines the conditions under which the physical system appears. Given two partial systems \(A\) and \(B\) forming a total system described by the state vector \(|\psi\rangle\), then we can ask the question: can we ascribe mutually independent existence (reality) to the partial systems \(A\) and \(B\) viewed separately, even if the partial systems are spacelike separated from each other at a particular time? Can the system \(B\) be directly influenced by measurements taken at \(A\)? The nonlocal Einstein, Podolsky and Rosen (EPR) paradox [8] taken in the context of these questions forces us to relinquish one of the following two statements [4]:

1. The state vector \(|\psi\rangle\) provides a complete description of the system.

2. The real states of spacelike separated systems are independent of each other.

If we choose to regard \(|\psi\rangle\) as a statistical ensemble of systems, thereby relinquishing (1), and giving up on the orthodox interpretation of quantum theory as a complete description of nature, then we will not find any contradiction with the standard quantum theory. The Copenhagen interpretation of quantum mechanics stipulates that quantum states do not exist before or after they have “jumped” into an eigenstate associated with a real eigennumber, following the discontinuous wave function reduction triggered by a macroscopic measuring device.
In the following, we shall develop a proposal to reconcile classical physics, relativity and QM. To this end, we shall formulate a wave function reduction, which involves an absorber mechanism associated with the interaction of the measuring apparatus with the past and future light cones. The formalism provides an explanation for nonlocal quantum measurements and a resolution of the EPR paradox that does not violate special relativity, i.e. it does not introduce superluminal velocities. Our theory of quantum measurements is related to the transactional interpretation of QM proposed by Cramer [8,9], who describes quantum mechanical wave functions as real waves physically present in space. The transaction is a quantum event describing an exchange of advanced and retarded waves; it is explicitly nonlocal and consistent with Bell’s inequality, but is relativistically invariant and causal.

II. THE ABSORBER THEORY OF RADIATION

In classical electromagnetism, certain solutions of Maxwell’s equations are discarded for empirical reasons. The field equations are time-symmetrical and show no preference whatever between the standard retarded fields of experience, which diverge at a finite speed from the source charges, and the corresponding advanced fields, which converge on source charges with the same speed. The advanced fields are discarded in favor of the retarded fields purely due to our selected experience. However, as we shall now argue, this experience may no longer be universally acceptable due to the demonstrated nonlocal nature of QM, and the need to conform with the laws of special relativity.

The process of choosing retarded versus advanced fields, which apart from our considerations of physical experience is an arbitrary choice, does not appear universal. Dirac [11] discovered that in order to describe the empirically well-established formula for radiation damping in terms of a covariant electromagnetic field, it is necessary to use both the retarded and advanced solutions of Maxwell’s equations.

We shall postulate that the advanced fields play a fundamental part in the interpretation of quantum measurements in QM. The role of advanced fields was considered in detail in two papers by Wheeler and Feynman [12,13], in which they proposed the absorber theory of radiation. They postulated that the time-symmetric solution, corresponding to one half the retarded plus one half the advanced fields, was the fundamental solution and that the arrow of time is generated by an absorber mechanism associated with all the other charges in the universe. In our construction of a wave function reduction mechanism, we shall use some of the ideas of this theory but expressed in a different way. We shall focus on the interrelation of the retarded and advanced fields and the absorber mechanism as correlated with the experimental apparatus and the wave function reduction. We shall find that in our interpretation, the wave function reduction is constrained by cosmological models. Indeed, in conjunction with the radiation absorber theory of wave function reduction, the electrodynamical arrow of time is a consequence of certain cosmological models and is incompatible with others. The connection between absorber radiation theories and cosmology was first pointed out by Hogarth [14], and was developed further by Hoyle and Narlikar [15,16], and Davies [17,18].

Consider a solution of Maxwell’s equations with particle \( i \) as the only source, and which has only retarded outgoing waves admitted. We denote this solution by \( F_{\mu\nu}^{(i)} \), while the corresponding advanced solution, which admits only incoming waves, is designated \( F_{\mu\nu}^{(i)} \). From Maxwell’s linear equations it follows that for any arbitrary real number \( a \), \( F_{\mu\nu} \) is a solution where

\[
F_{\mu\nu}^{(i)} = aF_{\mu\nu}^{(i)} + (1-a)F_{\mu\nu}^{(i)},
\]

Moreover, any superposition of the fields \( F_{\mu\nu}^{(i)}, F_{\mu\nu}^{(j)} \) etc., is a solution of Maxwell’s equations that takes into account all the sources. The familiar choice based on experience is that \( a \) is equal to unity. The assumption of Wheeler and Feynman, based on the action at a distance theory of electromagnetism formulated by Schwarzschild, Tetrode and Fokker [19,21] is that \( a \) equals 1/2, so that the physically significant solution of Maxwell’s equations is the time-symmetric field

\[
F_{\mu\nu}^{(i)} = \frac{1}{2}(F_{\mu\nu}^{(i)} + F_{\mu\nu}^{(i)}).
\]

They also asserted that there is no self-action of a particle on itself:

\[
F_{\mu\nu}(i) = \sum_{j \neq i} F_{\mu\nu}^{(j)},
\]

where \( F_{\mu\nu}(i) \) denotes the total field acting on particle \( i \).

Dirac [11] showed that the force of radiative reaction is relativistically deduced by means of the time-symmetric field described by
\[ F^{(i)}_{\mu\nu}^{\text{react}} = \frac{1}{2}(F^{(i)}_{\mu\nu}^{\text{ret}} - F^{(i)}_{\mu\nu}^{\text{adv}}), \]  

which acts only on particle \( i \). Observational classical electrodynamics is governed by the field

\[ F_{\mu\nu}(i) = \sum_{j \neq i} F^{(j)}_{\mu\nu}^{\text{ret}} + \frac{1}{2}(F^{(i)}_{\mu\nu}^{\text{ret}} - F^{(i)}_{\mu\nu}^{\text{adv}}). \]  

The time-symmetry of this equation also allows the equation to hold:

\[ F_{\mu\nu}(i) = \sum_{j \neq i} F^{(j)}_{\mu\nu}^{\text{adv}} - \frac{1}{2}(F^{(i)}_{\mu\nu}^{\text{ret}} - F^{(i)}_{\mu\nu}^{\text{adv}}). \]

An absorber theory of radiation will act in the following way. At each point \( O \) on a world line, sources of \( F(i) \) lie on the null cone with apex at \( O \). The null cone and the system of particles on it are called the absorber of \( i \) at \( O \). It was shown by Hogarth \[14\] that a static, infinite Minkowski universe, as used by Wheeler and Feynman, leads to an indeterminate solution for the absorber, whereas a non-static (expanding or contracting) universe can determine whether the absorber is ideal or non-ideal. Here, ideal or non-ideal refer to whether the absorber can produce either of the solutions, Eqs. \( (5) \) or \( (6) \). The conformal invariance of Maxwell’s field equations can provide solutions of the inhomogeneous equations in a conformally flat universe. Combined with the fact that essentially all homogeneous and isotropic solutions of Einstein’s gravitational equations are conformally flat, we can test the different absorber models for certain classes of conformally flat expanding universes, including the Einstein-deSitter universe.

The interactions of the particles in the absorber can be described by assuming propagation through a medium of refractive index. Wheeler and Feynman used the well-known formula for the refractive index of a medium consisting of unbound charged particles. Hoyle and Narlikar \[15\] used the theory in which the imaginary part of the refractive index arises from radiative reaction, rather than collisional damping. Davies used thermodynamic considerations and refractive index theory without complicated calculations involving Riemannian geometry to derive general conditions for determining the opaqueness of cosmological models \[16,18\].

### III. WAVE FUNCTION REDUCTION AND THE REVERSAL OF THE ARROW OF TIME

Both retarded and advanced wave solutions are consistent with relativity, in the sense that the propagation of signals occurs at the speed of light and there is no superluminal communication of information. This is in agreement with the observational result that no speed of propagation exceeds the speed of light in vacuum. On the other hand, the idea of causality is based on the notion that some temporal event occurs before another event. If \( O \) is a point in spacetime, a light cone determined by \( ds^2 = 0 \) belongs to it, where \( ds^2 \) is the square of the local Minkowski spacetime distance. We draw a timelike world line through \( O \) and on this line observe the close spacetime points \( X \) and \( Y \), separated by \( O \). If it is possible to send a signal from \( Y \) to \( X \), but not from \( X \) to \( Y \), then the one-sided, asymmetrical character of time is secured, and there exists no free choice for the direction of the arrow of time.

Because the wave equations for fields (including the electromagnetic fields) do not have an asymmetrical sense of the direction of time built into them, we can entertain at the quantum mechanical level the possibility that the arrow of time is reversed for a short duration during a measurement of the properties of a particle, due to the reduction of the wave function \( \psi \) describing an entangled system of particles including the particle being measured. The choice of the advanced Green function is then subject only to boundary conditions chosen on the basis of observations. We shall argue that nonlocal quantum theory observations do select the advanced Green function boundary conditions during the collapse of the wave function.

We shall adopt the usual approach to QM theory that the Schrödinger equation

\[ i\hbar \frac{\partial \psi}{\partial t} = H\psi \]  

describes a purely deterministic, unitary evolution of the wave function \( \psi \). This also holds true for the fields in a relativistic quantum field theory. However, this abruptly ceases to be true when we perform a measurement of the wave function \( \psi \), which destroys the coherence of the state vector by triggering a reduction of the the wave function. We shall make the following postulates:

1. Before and after a measurement by a macroscopic apparatus, with a corresponding reduction of the wave function, the electromagnetic field is described by the retarded solution given by Eq. \((6)\).
2. For purely local measurements of the particle the electromagnetic field is given by the retarded solution, Eq. (9), and there is no effect felt by the particle at \( O \) from an absorber stimulus either in the future or in the past light cone. Thus, for local measurements of a particle and its properties, the measuring apparatus does not interact with the absorber mechanism in the past and future light cones.

3. During a nonlocal measurement in the short time interval \( t_1 < t < t_2 \), the wave function reduction is correlated with an interaction of the particle being measured, situated at the apex \( O \) of the light cone, with the absorber mechanism in the past or future light cones. This has the effect of time reversing the field at the particle into the solution given by Eq. (9).

We must now guarantee that transmission of information takes place between the two entangled systems \( A \) and \( B \) which are spacelike separated, without invoking superluminal speeds. No transmission of information can occur directly between \( A \) and \( B \) at the speed of light, since they are spacelike separated events. We postulate that there exists an observer \( C \) at the apex \( O \) of a light cone that contains both the systems \( A \) and \( B \), such that the 45° angle cones formed at \( C \) intercept \( A \) and \( B \) in the future. This permits transmission of light signals between \( A, B \) and \( C \). When a measurement is performed, say, at \( A \), the time reversal of the electromagnetic field allows an advanced field to transmit the results of the measurement back in time to \( C \), which in turn transmits the information by a retarded field along its light cone to \( B \). In this way it would appear that information about the measurement is relayed instantaneously to \( B \). This is accomplished without violating special relativity, i.e. the signal communications between \( A \) and \( C \), and subsequently \( C \) and \( B \) occur at the speed of light in vacuum and local Lorentz invariance is maintained. The mechanism of using advanced fields to propagate information between \( A \) (or \( B \)) and \( C \) is also the basis of Cramer’s transactional interpretation of QM [9,10].

We have only used electromagnetic wave light signals to transmit the information between observers \( A, B \) and \( C \). We assume that the physical collapse of the wave function, triggered by a quantum nonlocal measurement, transmits information along the past and future light cones by means of photons, i.e. photon emission by the measuring device at the instant of measurement conveys the information about the physical state of observer \( A \) (or \( B \)) backwards in time to \( C \), where a device transmits information by means of photons to \( B \) (or \( A \)).

Cramer generalized the transactional mechanism to quantum wave functions for massive particles such as electrons or protons. The non-relativistic Schrödinger equation

\[
-\frac{\hbar^2}{2m} \nabla^2 \psi = i\hbar \frac{\partial \psi}{\partial t},
\]

where \( m \) is the mass of the particle is a first order equation in the time variable. Therefore, it does not possess advanced solutions. To overcome this problem, Cramer proposed using the reduction of a suitable relativistic wave equation to two distinct Schrödinger equations by taking the non-relativistic limit. The two resulting equations would be the complex conjugate or time reverse of one another. The time reversed Schrödinger equation,

\[
-\frac{\hbar^2}{2m} \nabla^2 \psi = -i\hbar \frac{\partial \psi}{\partial t},
\]

only has advanced solutions, and possesses negative energy eigenvalues.

One difficulty with this approach is the commonly held belief that correct quantum mechanical wave equations should be first order with respect to the time variable, so as to uniquely determine the state and the time evolution of the system. Moreover, the concept of retarded and advanced waves is generically a classical notion. Feynman’s propagator \( D_F \) in quantum electrodynamics is completely symmetric under time reversal, so that if we interchange the emitter and absorber labels of a photon under time reversal, then the resulting physical situation occurs with equal probability compared to the old one. However, in classical physics radiation phenomena involve large numbers of photons traveling in different directions in a correlated manner, so that the absence of converging waves can be attributed to the absence of correlation between different parts of the universe.

By making the transmission of “nonlocal” information with only photons, we avoid the problem of quantum wave transmission by advanced wave solutions associated with massive particles. Our transmission of information for nonlocal measurements is purely an electromagnetic wave phenomenon associated with the quantum measurement device and the collapse of the wave function.

**IV. ABSORBER MECHANISM AND COSMOLOGY**

The criterion for the existence of a complete absorber is based on arguments of the attenuation of the retarded and advanced fields as \( r \) increases to infinity. We shall assume that at the instant the apparatus makes a nonlocal
measurement and the wave function collapses, the particle $i$ being measured experiences a force of acceleration due to the time symmetric field $\frac{1}{2}(F_{\text{ret} \mu \nu}^{(i)} + F_{\text{adv} \mu \nu}^{(i)})$. The field from $i$ will produce an acceleration of other particles, which will in turn radiate retarded and advanced fields, so that the total field is

$$\sum F_{\mu \nu} = \frac{1}{2}(F_{\text{ret} \mu \nu}^{(i)} + F_{\text{adv} \mu \nu}^{(i)}) + \frac{1}{2} \sum_{j \neq i} (F_{\text{ret} \mu \nu}^{(j)} + F_{\text{adv} \mu \nu}^{(j)}), \quad (11)$$

the summation on $j$ corresponding to the effect of the past and future light cones in the rest of the universe. The expansion of the universe will break the symmetry and produce an arrow of time during the reduction of the nonlocal, entangled wave function.

The deterministic evolution of the fields is governed by Maxwell’s equations with the familiar retarded solution for the fields acting on the charges and also when a local measurement is made of the particles and their properties. When a measurement is made of a particle $a$ entangled with another particle $b$ at a spacelike separation $|x - y|$ from $a$, then the measuring device detects the total field $\sum_{i} F$ and the instantaneous collapse of the wave function converts this field into the total advanced field

$$\sum F_{\text{adv} \mu \nu} = F_{\text{adv} \mu \nu}^{(i)} + \sum_{j \neq i} F_{\text{adv} \mu \nu}^{(j)}. \quad (12)$$

If we wish to believe in our absorber wave function collapse scenario, then we have the observational cosmological constraint on the model that the universe should be transparent to light on the future null cone. This is a severe constraint on a cosmological model. It was first shown by Hogarth [14] that if the universe is an ideal (perfect) or quasi-ideal absorber along the future light cone, and is non-ideal along the past light cone, then the effective electromagnetic field acting on a particle at the apex of the light cone is a retarded field: $F_{\text{rad}}^{(i)} = 1/2(F_{\text{ret}}^{(i)} - F_{\text{adv}}^{(i)}), F_{\text{eff}}^{(i)} = F_{\text{ret}}^{(i)}$ observed in local measurements. For the opposite case, the field is the advanced electromagnetic field plus radiation: $F_{\text{rad}}^{(i)} = -1/2(F_{\text{ret}}^{(i)} - F_{\text{adv}}^{(i)}), F_{\text{eff}}^{(i)} = F_{\text{adv}}^{(i)}$.

By considering thermodynamic properties of cosmological models, Davies [17,18] has derived concise conditions to be satisfied by opaque and transparent universes. Let us consider a single photon. The probability that a photon will be absorbed in time $dt$ while passing through objects of density $\rho$ and cross section $\sigma$ is

$$1 - \exp(-\rho \sigma dt) \sim \rho \sigma dt. \quad (13)$$

If the integral

$$\int_{-\infty}^{\infty} \rho \sigma dt = \infty, \quad (14)$$

then the probability is unity. For constant $\sigma$, Eq.(14) has the limiting case $\rho \propto 1/t$ or $R \propto t^{1/3}$, where $R = R(t)$ denotes the scale factor in homogeneous isotropic cosmological solutions of Einstein’s field equations [22]. For steady-state cosmology both $\rho$ and $\sigma$ are constant, so Eq.(14) is satisfied.

The mean cross section for photon absorption by an electron is

$$\sigma = \frac{A \rho}{T_{i}^{1/2} \omega^3} [1 - \exp(-\omega/kT_{i})], \quad (15)$$

where $T_{i}$ is the ion temperature, $\rho$ is the total heavy particle density and $A$ is a numerical factor. We now use that $\rho, T_{i}$ and $\omega$ are proportional to $1/R^3, t^{2/3}/R^{8/3}$ and $1/R$, respectively, which yields $\sigma \propto R^{4/3}/t^{1/3}$. The density of electrons falls off as $1/R^3$, so complete absorption results in

$$\int_{0}^{\infty} R^{-3} t^{-1/3} R^{4/3} dt = \infty \quad (16)$$

giving the limiting case $R \propto t^{2/5}$.

Thus, the general conclusion is that all expanding matter conserving cosmological models do not permit complete (ideal) absorption along the future light cone. However, they lead to complete absorption along the past light cone due to the big-bang singularity and the increasing density of matter and radiation as the singularity is approached. For recontracting models which evolve to high density models in the future as the universe recollapses to a singularity, complete absorption can occur, leading to consistency with a retarded solution instead of an advanced solution. However, such models do not lead to favorable observational density values. Steady-state models also yield consistent
We have constructed a theory of quantum mechanics that incorporates a wave function reduction process based on the absorber theory of electrodynamics. During the short time that an apparatus detects the properties of a particle associated with an entangled state vector, the electrodynamic arrow of time is reversed, whereby a premonitory signal is activated between the location of a measurement and an observer situated in the past light cone of the measurement.
event, who in turn communicates the results of the measurement to the spacelike separated event associated with the other state of the entangled system. The exchange of information occurs at the speed of light. Measurements of local quantum observables do not involve such a reversal of the electrodynamic arrow of time, so that only the retarded field influences the acceleration of particles. Moreover, when no measurements are performed, the classical electromagnetic fields obey the standard retarded field solutions, the wave function satisfies the deterministic, unitary Schrödinger equation, and quantum fields obey deterministic relativistic field equations. All the paradoxes associated with EPR experiments are resolved, for the anti-causal information exchange between spacelike separated systems makes the wave function reduction consistent with special relativity and quantum mechanics.

Expanding universes in conventional cosmology reverse the direction of the arrow of time during the wave function collapse, selecting the advanced electromagnetic field boundary conditions for nonlocal measurements of systems that are spacelike separated. Thus, standard expanding FRW universes with zero, positive or negative spatial curvature are consistent with the theory and with current observations in cosmology, which prefer a big-bang scenario with $\Omega_{\text{crit}} = \rho/\rho_{\text{crit}} \leq 1$. It is interesting that the absorber model of wave function reduction is restricted by the cosmological observations due to its Machian-type property of depending on all the charges in the entire universe.

The absorber wave function reduction theory with its associated reversal of the arrow of time is expected to be a strictly quantum phenomenon like the exclusion principle or quantum spin; there is not anticipated to be a classical analog of this process, so that macroscopic premonitory signals propagated at the speed of light are forbidden to exist. Thus, there is no conflict of this theory with the second law of thermodynamics, i.e. the arrow of time created by the increase of entropy at the macroscopic level will not conflict with the quantum reversal of the arrow of time. However, the theory does open the possibility of quantum premonitory exchange of information which could have fundamental importance for quantum computer devices or other quantum devices; the EPR phenomenon has already produced interesting proposals for quantum cryptography [25].

The idea of an advanced electromagnetical wave exchange of information would suggest that at the quantum level of the world there is no free will, since all events in the quantum world are preordained for spacelike separated entangled systems. This issue arises because we associate a physical reality with the spacelike correlated systems and there is a finite speed of communication of information between such systems that is not in conflict with relativity.

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