N-body simulations of warped galaxies

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Abstract.
Two methods generating warped galaxies with N-body simulations are presented. One uses an external potential as a disturber while the other is based on material accretion. The results of both methods are compared. A particular attention is given on the shape of the line of node (LON).

1. Introduction
The understanding of warped galaxies has been enriched those last years by ideas such as galactic infall (Ostriker & Binney 1989; Jiang & Binney 1999) or dynamical friction between the halo and the disk (Debattista & Sellwood 1999). In spite of the fact that the evidences for non-self-gravitating HI disks are actually weak (Arnaboldi et al. 1997; Becquaert & Combes 1997), most of works on the subject are built on the assumption of a disk embedded in a fat spheroidal or triaxial halo.

Through N-body simulations, this work presents two methods of getting warped self-gravitating galactic disks. The first one imports inclined angular momentum through material accretion and the second one uses an outer warped potential.

2. Models and simulations
The N-body disk is based on a superposition of three Miyamoto-Nagai potentials:

\[ \Phi(R, \phi, z) = -\sum_{i=1}^{3} \frac{GM_i}{\sqrt{R^2 + (a_i + \sqrt{z^2 + b_i^2})^2}} \]  

(1)

with the components representing a bulge, a visible disk and a gas disk.

2.1. External potential
In order to warp the galaxy, an initially flat disk is embedded in a non-rotating warped potential. This outer potential whith a mass equivalent to the disk is obtained by replacing \( z \) by \( z - wR^2 \cos \phi \) in the disk potential of equation 1. \( w \) is a parameter determining the amplitude of the warp. To drive adiabatically the disk, the outer potential is flat at the beginning of the simulation and warps
slowly during the perturbation time. Afterwards, the potential is removed and the disk evolves freely.

After a few tens of Myr, the inner part of the disk tilts while the outer remains flat. At $t \approx 750$ Myr the whole disk has warped in order to follow the outer potential and the LON is observable as a leading spiral. The maximum amplitude of the warp at $R = 30$ kpc is about 20 kpc. After the perturbation, the deformation of the isolated disk decreases from the center to the outside and the LON becomes straight. This last step needs about 1 Gyr.

2.2. Accretion

In this simulation, the warp is generated by material accretion. This idea has been suggested by Ostriker & Binney (1989) and tested by Jiang & Binney (1999). The originality of our simulation is the absence of halo. The infall of material is obtained by injecting particles in an inclined torus. Its velocity corresponds to the velocity of the rotation curve at an equivalent radius. The total injected mass represent 100% of the mass of the initial galactic disk.

The disk needs about 1 Gyr in order to react to the accretion and the deformation will need one more Gyr to propagate from the outside to the center. The maximum amplitude of the warp at $R = 30$ kpc is about 10 kpc. A straight retrograde rotating LON is associated to the warp. Once the accretion is ended the disk needs about 500 Myr to become flat again.

3. Discussion

The perturbation type determines strongly the evolution and shape of the warped disk. In the first simulation, the deformation occurs from the inner regions to the outer whereas it occurs from the outer to the inner in the second one. This has two observed effects:

The shape of the LON is significantly different. A leading LON is generated in presence of an outer potential, which can be assimilated to a warped heavy DM disk. A similar result has been found by Debattista & Sellwood (1999) in the case of dynamical friction between a spherical misaligned rotating halo and a disk. In contrary, a straight rotating LON is generated with accretion.

The birth of the warp is also quite different. In case of our outer potential, the inner regions feel its gravity faster and tip before the outer ones. A large abrupt warp can be quickly observed, even if the outer regions lie always on the initial plane. This effect is not present in case of accretion, where the inner regions remain flat while the outer ones tip slowly.

References

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