Polarization selective phase-change nanomodulator

Kannatassen Appavoo1,2 & Richard F. Haglund Jr.1,3

1Interdisciplinary Materials Science, Vanderbilt University, Nashville, TN 37235-0106, USA, 2Center for Functional Nanomaterials, Brookhaven National Laboratory, Upton, New York 11973, USA, 3Department of Physics and Astronomy, Vanderbilt University, Nashville, TN 37235-1807, USA.

Manipulating optical signals below the diffraction limit is crucial for next-generation data-storage and telecommunication technologies. Although controlling the flow of light around nanoscale waveguides was achieved over a decade ago, modulating optical signals at terahertz frequencies within nanoscale volumes remains a challenge. Since the physics underlying any modulator relies on changes in dielectric properties, the incorporation of strongly electron-correlated materials (SECMs) has been proposed because they can exhibit orders of magnitude changes in electrical and optical properties with modest thermal, electrical or optical trigger signals. Here we demonstrate a hybrid nanomodulator of deep sub-wavelength dimensions with an active volume of only 0.002 \( \text{m}^3 \) by spatially confining light on the nanometre length scale using a plasmonic nanostructure while simultaneously controlling the reactive near-field environment at its optical focus with a single, precisely positioned SECM nanostructure. Since the nanomodulator functionality hinges on this near-field electromagnetic interaction, the modulation is also selectively responsive to polarization. This architecture suggests one path for designing reconfigurable optoelectronic building blocks with responses that can be tailored with exquisite precision by varying size, geometry, and the intrinsic materials properties of the hybrid elements.

As the transistor reaches its fundamental size and speed limits, new elementary switching structures based on optoelectronics are being developed for next-generation computers1–2. Plasmonics has the potential to create the requisite optoelectronic building blocks by integrating electronics and photonics on a single chip, enabling the exchange of data at larger bandwidth and at ultrafast speeds3. Although light can be squeezed and manipulated at sub-wavelength scale, recently reaching confinement dimensions of \( \lambda/165 \) in quantum metallic nanostructures4 and \( \lambda/200 \) for point defects in graphene5, dynamically modulating the intensity or phase of optical signals at these dimensions at THz frequencies and at a maximal energy cost per bit of 100 fJ remains a challenge6.

Active modulation of plasmon optical response has been demonstrated by optical excitation of a thin layer of quantum dots7, the electro-optic effect in barium titanate thin films8, optically induced refractive index changes in silicon9 and recently, using transparent conductive oxides10. In optical modulation of the surface plasmon polariton by quantum dots, for example, energy costs of less than 1 fJ per bit were derived from a thin-film geometry, but the speed would be limited to sub-GHz range because of the spontaneous lifetime of the optical transition. The refractive-index change in silicon is ultrafast, but the modulation depth is small because of the small index change (\( \Delta n = 0.005 \)) available from the electro-optic effect in silicon. And in all these operational schemes, it appears that plausible thin-film devices would have areas of several square micrometres.

Here we demonstrate active modulation of optical-frequency plasmonic resonance peaks by about 12 nm in a device whose footprint is less than 0.1 \( \mu \text{m}^2 \) by using a reversible phase transition (PT) in a prototypical strongly electron-correlated material (SECM), vanadium dioxide (VO2). Vanadium dioxide undergoes a metal-to-insulator (MIT) phase transition near 68 \( \mu \text{C} \) with a concomitant atomic rearrangement from tetragonal to rutile that leads to dramatic changes in optical and electrical properties. The phase transformation can be driven optically on time scales less than 100 fs11–13 and as recently demonstrated by Wall et al.14,15, the ultrafast PT can be modulated optically at low fluences by femtosecond excitation of specific VO2 vibrational modes. Moreover, VO2 nanostructures exhibit size-dependent switching as well as a dipolar plasmon response in the near-IR when the nanostructures metallize16,17.

As schematically represented in Fig. 1a, the intrinsic plasmonic response of VO2 nanostructures is strongly damped17, limiting their utility when deployed as the sole switchable component in a nanostructure. However, the large change in dielectric function – of order unity18 – between the “on” and “off” states in a single VO2 nanodisk...
placed in the optical near field of a metal nanodisk enables effective modulation of the plasmon resonance (Fig. 1b vs. 1c). This modular nanoarchitecture – relying solely on the near-field electromagnetic interaction between the metallic and SECM nanostructure – enables various functions, as illustrated in Fig. 1d. For example, spectral response from the visible to the mid-IR is obtained by varying the size and shape of the metallic nanostructure. The intrinsic phase-changing properties of the VO₂ can also be harnessed to create bistable switches or colour nanorouters by exploiting thermally, optically or electrically driven, size-dependent VO₂ switching.

This design flexibility is achieved by controlling the size of the VO₂ nanodisks, dictating the density of defects that nucleate the phase transition and thus hysteresis width (or operational range) of the nanomodulator (Fig. 1d–I). Nanomodulators operating at THz frequencies can be envisioned in the non-equilibrium limit (Fig. 1d-II). Furthermore, since the hybrid nanomodulator functions via the electromagnetic near-field interaction, selectivity is achieved by polarization as well (Fig. 1d-III), as will be experimentally and computationally demonstrated below.

The prototype hybrid nanomodulator was fabricated using a three-stage electron-beam lithographic method (EBL) illustrated in Fig. 2 (see Methods). Such a fabrication protocol, to the best of our knowledge, has been successfully achieved only by Liu et al. for gas sensing using a gold/palladium nanoarchitecture. Here, lithographically fabricated gold nanomarkers were used to position a gold nanodisk on an indium-tin-oxide (ITO)-covered glass substrate in a lattice configuration. Using the same nanomarkers, a single VO₂ nanodisk was subsequently positioned adjacent to its neighbouring gold nanodisk. The sample was then annealed to render the VO₂ nanostructure crystalline (insulating-M1) and phase transforming. As shown in the scanning electron micrograph (SEM, Fig. 2a), the differential contrast of the two nanostructures clearly shows the efficacy of the design. To demonstrate that the nanomodulator operates solely by near-field electromagnetic interactions, we fabricated ten similar patterns on that same substrate with varying separations between the gold and VO₂ nanodisks (Fig. 2c). Using this protocol, we achieved a minimum edge-to-edge nanostructure separation of about 14 nm as confirmed by SEM in Figs. 2a, c.
Following fabrication, white-light broadband extinction was measured on each nanomodulator array both when the VO$_2$ nanostructure was in the insulating (blue, room temperature) and metallic (red, 95°C) states (Fig. 3a, see Methods), with the light polarized parallel to the long axis of the nanomodulator. When the VO$_2$ nanodisk is located at a distance of 180 nm from the gold nanodisk, no modulation of the plasmonic extinction spectrum is observed. However, at an edge-to-edge separation of about 34 nm, the nanomodulator shows the first signs of spectral tuning. As the VO$_2$ nanostructure approaches the near-field of the gold nanodisk, the wavelength shift increases. At the minimum separation that could be achieved with the available EBL tool, a blue shift in the modulation of about 12 nm between the spectral peaks of the hot-metallic and cold-insulating states is obtained. This is better than or comparable to previous studies using thin films as active switching media.

Consequently, this proof-of-concept experiment demonstrates that a single nanostructure of strongly electron-correlated material carefully positioned with respect to the plasmonic nanostructure can effectively modulate optical signals in a nanoscale volume. The high contrast in dielectric function between the metallic and insulating states of the VO$_2$ makes it possible to achieve effective modulation in substantially smaller volumes than other material systems; in the present case, the active volume is less than 0.002 μm$^3$ compared to 5–100 μm$^3$ for conventional electro-optic modulators. If triggered optically, such a hybrid phase-change/plasmonic nanostructure has been shown to require at most 250 μJ/cm$^2$ or control surface plasmon-polariton propagation at an aluminium/silica interface (∼10000 μJ/cm$^2$). Given that the area of the modulator is of order 0.1 μm$^2$, a switching threshold of 250 μJ/cm$^2$ corresponds to an energy cost of ∼20 fJ per bit for a wavelength shift of 1 nm, assuming that the wavelength shift is to zeroth order linearly proportional to the absorbed energy. This is well within the canonical energy cost per bit for a photonic switching system. We note that when light is polarized along the short axis ($\mathbf{E}_\perp$), no modulation is observed experimentally.

To estimate the maximum possible modulation in this geometry, three-dimensional full-field finite-difference, time-domain simulations with light polarized either parallel (\(\mathbf{E}_/\)) or perpendicular (\(\mathbf{E}_\perp\)) to the axis of the nanomodulator were carried out, with nanostructure edge-to-edge separation as small as 0 nm (“kissing”) nanostructures. As shown in Fig. 3b, there is no modulation for edge-to-edge distances from 180 nm to 40 nm. However, when the VO$_2$ nanostructure reaches the electromagnetic near-field of the Au nanodisk, a progressively stronger modulation of the plasmon resonance wavelength between the hot and cold VO$_2$ states is achieved. As seen in Fig. 3c, the simulated and the experimental extinction spectra agree over the range of fabricated nanostructures. Modulation begins only when the two nanostructures are less than approximately 35 nm apart; inspection of Fig. 3b shows that this is in the weak-coupling regime where minimal plasmon hybridization occurs and the VO$_2$ nanodisk can be effectively treated as a dielectric perturbation. While a systematic redshift as the insulating VO$_2$ approaches the Au-nanostructure, we observe a slight blueshift when the nanostructure is in its metallic state, consistent with an intrinsic plasmonic response in the near-IR for a plasmonic density of states smaller than that of the gold nanostructure. More importantly, the simulations show that a wavelength modulation approaching 50 nm is achievable for the “kissing” nanostructures; however, this must be taken as only an indication of an extreme value, since at this spacing other mechanisms come into play, such as quantum tunnelling and electronic spill-out effects.

Experimental tests of the polarization dependence of the hybrid nanomodulator showed no sign of spectral tuning when light was polarized perpendicular to its axis. Since a separation of 14 nm was the smallest achieved experimentally, we used simulations to understand how polarization affects the behavior of the nanomodulator for smaller separations. Comparing the near-field electromagnetic profiles for parallel (\(\mathbf{E}_/\)) and perpendicular (\(\mathbf{E}_\perp\)) polarization (Fig. 3b), it is clear that the EM fields interact strongly with the VO$_2$ nanostructure only for parallel polarization. The field distortion created by the VO$_2$ nanostructure within the lobe of the Au dipolar resonance therefore dictates the overall modulation effect of the device. This implies that nano-focusing using more complex architectures like...
nanopyramids is not necessary as even with a relatively broader “spillout” of the plasmon near-field from the gold nanostructure, selective modulation can be achieved by polarization. This can be visualized in Fig. 4b (right) where the symmetric distortion of the dipole field with respect to the entire nanomodulator results in no spectral tuning, albeit with a slight broadening in linewidth, requiring a near-perfect positioning of a phase-changing nanostructure to maximize modulation (Fig. 3b).

Figure 3 | Experimental and simulated optical spectra of nanomodulators. (a), Experimental extinction spectra for the 172, 160, 129, 122, 102, 72, 66, 54, 34, 14 nm interparticle separation arrays when the vanadium dioxide nanostructures are in the metallic (red) or insulating (blue) state. (b), Simulated extinction spectra of the hybrid nanomodulator for interparticle separation of 180, 140, 120, 80, 60, 40, 20, 14, 10, 8, 6, 4, 2, 0 nm in metallic (red) and insulating (blue) state. The orange-shaded plots in each figure represent the simulated optical response of the gold lattices that match closely both with the experimental results and with the cases where the VO2 nanostructures are not in the near-field region of the plasmonic nanostructure. Spectral modulation is visible when edge-to-edge separation is less than 35 nm (c), Summary of spectral shifts for both experimental and simulated results. The error bar for the interparticle separation was estimated to be ± 4 nm while the peak determination is ± 2 nm. We note that all measurements (experimental and simulated) shown here are for light polarized along the long axis ($\vec{E}_l$).
Although the hybrid Au-VO₂ nanomodulator cannot be compared directly to conventional electro-optic modulators, its critical parameters – modulation speed, extinction ratio, bandwidth, power consumption, and device footprint – still measure up well. With a single nanostructure having an active volume of only 0.002 mm³, the modulation depth for this hybrid nanomodulator is ~1.3 dB, calculated by comparing the transmission intensity change at the peak plasmon resonance when the VO₂ nanostructure is metallic. However, increasing the modulation depth is straightforward by implementing strategies such as optimization of the size and shape of the VO₂ nanostructure (for example a split-ring circular structure) or implementing this near-field design in well-established photonic devices such as ring resonators⁹ or interferometers⁹. We note that our nanomodulator was not optimized at telecommunication wavelength but operation bandwidths greater than 100 nm can be realized with modulation speed of 1 GHz, while shifting the operational bandwidth can be achieved by varying the size of the plasmonic nanostructure. Although our polarization-selective hybrid nanomodulator has been demonstrated with two nanodisks of similar size, other geometries such as nanorods¹¹ to decrease the plasmonic line-lengths but operation bandwidths greater than 100 nm can be readily designed and fabricated. Additionally, even though this demonstration experimentally focused on modulation of a single localized surface-plasmon resonance, the application to modulating information-carrying chains of nanoparticles would appear to be straightforward as well¹²⁻¹⁵.

Technologically, this simple nanoarchitecture provides substantial design flexibility, enabling other potential applications requiring minimal footprints, such as trapping¹⁷, beam steering and focusing, where dictating the optical phase of light is critical¹⁸. Moreover, the unique nonlinear optical properties of VO₂ allow for non-centrosymmetric nanomodulator designs to access second- and third-order nonlinearities¹⁹. This simple architecture also makes it possible to study the interactions of ensemble nanostructures in which one or more components has an intrinsically variable plasmonic density of states²⁰.

**Methods**

**Fabrication.** Gold markers were lithographically fabricated on ITO-coated glass substrate by means of (1) electron-beam lithography (EBL) on PMMA, spun at 2500 rpm and baked at 180°C for 90 sec; (2) e-beam evaporation of 20 nm of gold on the patterned substrate, measured by a quartz-crystal microbalance; (3) subsequent lift-off procedure using a Remover PG bath at 70°C for 20 min. The gold nanodisk lattices were fabricated using the same protocol, resulting in nanodisks of height 20 nm and unit cell of 600 nm, each plasmonic lattice having a footprint of 100 × 100 μm². Using the gold alignment markers, a third step lithography was performed by coating once more PMMA resist on the sample and writing the phase-changing lattice pattern at sub-20-nm level. Subsequently, the vanadium dioxide nanodisks were generated by (i) pulse laser ablation of a vanadium metal target (PLD: λ = 248 nm, 25 ns pulse duration, 3.86 J/cm² fluence, 10 Hz repetition rate and 10 mTorr of O₂ gas) to yield, after subsequent lift-off, nanodisks of amorphous, sub-stoichiometric vanadium dioxide (40 nm VOX₁₋₂ nominal thickness); (ii) thermal anneal of the sample (450°C, O₂ gas at 250 mTorr for 20 min) to render the NPs crystalline and stoichiometric.

**Optical measurements.** Extinction measurements were acquired as follows: (i) the 100×100 μm² plasmonic lattice of interest was positioned using micrometre drives under white light illumination from a tungsten lamp (90 μm spot size and polarized at 45° (perpendicular) or 135° (parallel) to the interacting dimer) while concurrently being imaged onto a CCD camera in the x-y plane; (ii) the focus (z-plane) was then visually adjusted by displacing the 5X microscope objective with a numerical aperture of 0.12; (iii) extinction at room temperature was then measured (integration time of 8 ms and average of 200) (iv) the sample was then heated to 95°C, well above the transition temperature of the VO₂ nanodisks, and extinction measurements for the various lattices were repeated. As flat-field, transmission spectra of the ITO-coated glass at room temperature and at 95°C were used when the VO₂ nanodisks were...
insulating and metallic respectively. We note that this is in contrast to using the insulating or metallic VO₂ film as the flat-field insulator or metallic respectively. We note that this is in contrast to using the insulating or metallic respectively. We note that this is in contrast to using the insulating or metallic respectively. We note that this is in contrast to using the insulating or metallic respectively.