The frequency dependence of scattering imprints on pulsar observations

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ABSTRACT

Observations of pulsars across the radio spectrum are revealing a dependence of the characteristic scattering time ($\tau$) on frequency, which is more complex than the simple power law with a theoretically predicted power-law index. In this paper, we investigate these effects using simulated pulsar data at frequencies below 300 MHz. We investigate different scattering mechanisms, namely isotropic and anisotropic scattering, by thin screens along the line of sight, and the particular frequency-dependent impact on pulsar profiles and scattering time-scales of each. We also consider how the screen shape, location and offset along the line of sight lead to specific observable effects. We evaluate how well forward fitting techniques perform in determining $\tau$. We investigate the systematic errors in $\tau$ associated with the use of an incorrect fitting method and with the determination of an off-pulse baseline. Our simulations provide examples of average pulse profiles at various frequencies. Using these, we compute spectra of $\tau$ and mean flux for different scattering setups. We identify setups that lead to deviations from the simple theoretical picture. This work provides a framework for interpretation of upcoming low-frequency data, both in terms of modelling the interstellar medium and understanding intrinsic emission properties of pulsars.

Key words: scattering – pulsars: general – ISM: structure.

1 INTRODUCTION

Observed radio pulses from pulsars pass through the ionized interstellar medium (ISM) before terrestrial detections are made. As the waves propagate through the ISM, inhomogeneities in the electron densities of dense, ionized regions cause the radio waves to scatter. This scattering is observed through, amongst other effects, the frequency-dependent temporal broadening of the received radio pulses. The temporal broadening typically leads to characteristic exponential tails in averaged pulsar profiles (e.g. Löhrer et al. 2001). Understanding the frequency dependence of the temporal broadening allows us to investigate the properties of the scattering media. Compensating for scattering effects can be used to infer the intrinsic radio emission characteristics of pulsars.

The scattering tails of the broadened profiles are routinely modelled with exponential functions that take the form $\tau^{-1}e^{-\tau/\tau}$, where $\tau$ is referred to as the characteristic scattering time. This functional form, as described in detail in Section 2.1, is based on the assumption of a single, thin scattering screen extending infinitely transverse to the line of sight to the pulsar. The screen is invoked as an approximation of the combined scattering that the pulse signal undergoes in the extended ISM along the line of sight (Scheuer 1968; Cronyn 1970; Williamson 1972, 1973). The characteristic scattering time $\tau$ depends on the properties of the scattering screen as well as its location along the line of sight (see equation 2) and it is maximized for a mid-way screen.

Apart from approximating the overall scattering geometry to a single scatterer, historic studies have also investigated distinct electron density distribution models for the screen. These include an isotropic Gaussian electron column density wavenumber spectrum, which leads to a temporal broadening dependence on frequency, $\tau \propto v^{-\beta}$, (Lang 1971; Lee & Jokipii 1976). Alternatively power-law wavenumber spectra, characterized by an inner and outer cut-off length scale, are adopted to model the plasma inhomogeneities. Most often, a Kolmogorov power spectrum, which describes the turbulence of a neutral gas, is preferred (Rickett 1977; Rickett, Coles & Bourgois 1984; Rickett 1990). A Kolmogorov spectrum will have a frequency dependence of $\tau \propto v^{-4.4}$, for wavenumbers lying between the inner and the outer cut-off length-scales.\textsuperscript{1}

The scale at which electron density inhomogeneities lead to strong diffractive scintillation and consequently the observed pulse

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\textsuperscript{1} Note that for power-law spectra, where $\tau \propto v^{-\beta}$, the fluctuation spectral index $\beta$ is often quoted instead of the spectral index $\alpha$. In the 2D case, for $\alpha \geq 4$, they are related through: $\alpha = 2\beta/(\beta - 2)$. A Kolmogorov spectrum with $\alpha = 4.4$ has $\beta = 11/3$. Also note that the symbols $\alpha$ and $\beta$ are sometimes used interchangeably.
broadening analysed in this paper, is set by the field coherence scale, \(s_0 = 1/(k\delta)\), with \(k\) the radio wavenumber (Rickett 1990). For strong diffractive scintillation, \(s_0 \ll L_\text{F}\), where \(L_\text{F} = (D/k)^{1/2}\), is the Fresnel scale (e.g. Smith & Thompson 1988. Refer to Fig. 1 for descriptions of \(\theta\) and \(D\)).

Estimations of \(\tau\) values from pulsar profile observations are done by either deconvolution of the folded pulse profile with a broadening function, such as a one-sided exponential function through e.g. the CLEAN-based algorithms (Bhat, Cordes & Chatterjee 2003), or by forward modelling which convoluses a template unscattered pulse profile (or alternatively an observed high-frequency pulse profile) with a parametrized broadening function and compares the model to the observed profile (e.g. Löhmer et al. 2001, 2004). It should be noted that using a high-frequency observation as a template will, for pulsars that exhibit intrinsic frequency evolution, introduce biases to measurements of scattering parameters. A deconvolution method does not suffer from this effect.

Alternatively \(\tau\) can be estimated by means of the scintillation bandwidth, \(\delta f\) (Lyne & Rickett 1968). The scattering process imparts frequency-dependent phase changes on the passing radio waves. An interference pattern will emerge for phase changes of the order \(2\pi\delta f\tau \approx 1\) rad. A measurement of \(\delta f\), over which the interference effects are observed in the dynamic spectrum, can therefore lead to an approximation of \(\tau\) (Cordes 1986).

Using the ISM transfer function implied by the thin screen model results in fits of \(\tau\) in observed data that follow the expected spectral dependence (Armstrong, Rickett & Spangler 1995). However, deviations from the theoretical spectral dependences have also been proposed and observed – both with spectral index \(\alpha\) (where \(\tau \propto \nu^{-\alpha}\)) smaller than predicted (Löhmer et al. 2001; Bhat et al. 2004; Lewandowski, Kowalinska & Kijak 2015) and larger than predicted (Rickett 1990; Lambert & Rickett 1999; Tuntsov, Bignall & Walker 2013). Different models will lead to different values for \(\alpha\), but the critical minimum for strong diffractive scattering, is \(\alpha = 4\) (Romani, Narayan & Blandford 1986).

Anisotropic scatterers, i.e. media for which the distribution of scattering angles exhibit directionality have been considered as causes for steeper spectra (Tuntsov et al. 2013). They were first invoked to explain organized patterns in the dynamic spectra of pulsar observations (Gupta, Rickett & Lyne 1994). Stinebring et al. (2001) found that complex patterns in dynamic spectra are often related to parabolic arcs in the secondary (power) spectra of pulsar observations, which can be explained by anisotropic scattering (Walker et al. 2004). There has also been some successes in associating anisotropy with observations of intra-day variability (Bignall et al. 2003; Tuntsov et al. 2013). Extreme scattering events as observed by Fiedler et al. (1987) and analysed by e.g. Walker (2007) have motivated works studying the impact of alternative scatterers, such as self-gravitating, au-sized ionized gas clouds. A 1D Kolmogorov power spectrum in electron density distribution is often used to model extreme anisotropy.

On the contrary, weaker dependences on frequency have been attributed to screens which have a limited physical size rather than being infinite transverse to the line of sight as most models assume (Cordes & Lazio 2001). Beyond the edges of such a screen, no radio waves coming from the pulsar will be bent back into the line of sight of the observer, leading to no further scatter broadening with increasing wavelength. Such configurations also lead to a loss in observable flux.

A physical inner scale of the scattering screen, that is smaller than the diffractive scattering scale, will restrict diffraction to a maximum scattering angle. This too will lead to less power observed at large angles, or alternatively flatter \(\tau\) spectra with lower power-law indices, at long wavelengths. In the case of an inner scale, however, flux is preserved (Rickett et al. 2009).

Until recently, characterizing the frequency dependence of temporal scatter broadening was done using pulsar observations taken at different receivers and telescopes at different epochs. With the construction of telescopes such as Low-Frequency Array (LOFAR; van Haarlem et al. 2013) and the Murchison Widefield Array (MWA; Tingay et al. 2012), we have the advantage of studying pulsars at low frequencies, where scattering is most pronounced, using broadband data (see Stappers et al. 2011 for an outline of LOFAR’s pulsar observing modes). This allows accurate measurements of \(\tau\) spectra and simultaneous studies of the flux spectrum and profile evolution of the pulsar. Analysis of broad-band data will therefore test our current understanding of pulsar scattering.

Recent works showing average pulse profiles in the LOFAR bands (30–90 MHz and 110–190 MHz) confirm the expectation of scatter broadened pulse shapes at these frequencies (Bilous et al. 2016; Pillia et al. 2016).

In this paper, we provide a framework based on simulations, for interpreting observations such as the ones presented in the LOFAR works. This can be used to gain insight on the size, location and nature of scattering screens.

The paper is split into sections as follows. In Section 2, we describe the general scattering setup, the simulated data and the techniques by which we fit the data to retrieve \(\tau\) spectra. The mathematical formalism for examples of the broadening functions for isotropic and anisotropic scattering is also presented in this section. In Section 3, we present the results of our experiments. These include example profile shapes as well as \(\tau\) and flux spectra for infinite and finite screens, respectively. We end with a discussion of our results and their relevance to broad-band observations in Section 4, followed by a short conclusion.

## 2 SIMULATED DATA AND FITTING TECHNIQUES

### 2.1 Scattering geometry

The scattering geometry considered in this work is shown in Fig. 1. The radio emission from a pulsar at a distance \(D\) from the observer...
Scattering imprints on pulsar observations

2.2.1 Simulated intrinsic pulse

The intrinsic pulse profile in our simulated data consists of a single Gaussian component with a chosen duty cycle and pulse period \((P)\). When necessary, multiple pulse periods are stitched together to form a pulse train. We investigate a typical slow pulsar with \(P = 1\) \(s\) and a duty cycle of \(2.5\) per cent and fast pulsar with \(P = 20\) ms with a duty cycle of \(10\) per cent.

The mean flux spectral index for all the simulations is chosen to be \(\gamma = 1.6\), where the mean flux is \(S_m = v^{-\gamma}\). We use \(\gamma\) as the flux spectral index and \(\alpha\) as the \(\tau\) spectral index to avoid confusion. The mean flux of the pulsar is normalized, such that it has a maximum value of unity for the lowest studied frequency. The use of a Gaussian template is motivated by the fact that a large fraction of pulsar profiles are relatively well approximated by this shape (e.g. Johnston & Weisberg 2006; Karastergiou & Johnston 2007.). The drawback is that a single Gaussian shape is not representative of all pulsar profiles, especially not of older, slower pulsars. In Section 3.3, we briefly study possible scattering impacts on a more complex profile shape based on the average profile of PSR B1237+25.

2.2.2 Simulated scattering screen

We consider nearby sources by choosing a standard distance of \(D = 3.0\) kpc from the observer to the pulsar with a mid-way scattering screen at \(D_s = 1.5\) kpc. We assume we can treat the scattering screen as a cold plasma for which the scattering angle \(a \propto v^{-2}\), with \(v\) the observing frequency (Lang 1971). For a distribution in scattering angles, the standard deviation will follow the relation \(\sigma_a \propto v^{-2}\). The scattering strength of the screen is fixed by choosing a proportionality constant in this relation. We use \(\sigma_a = 3\) mas at 1 GHz, such that \(\sigma_\theta = 1\) mas at distances \(D = 1\) kpc and \(D_s = 1/3\) kpc, as suggested in Cordes & Lazio (2001).

We present two analytic expressions for temporal broadening functions associated with isotropic and anisotropic scattering by an infinite thin screen, and thereafter discuss a ray-tracing experiment by which numerical broadening functions for alternative configurations can be built up.

Isotropic broadening functions. The temporal broadening function for a 2D scattering screen that scatters photons isotropically and can be described by a circularly symmetric Gaussian probability distribution in the scattering angle, takes the following form

\[
f_i = \tau^{-1} e^{-t/\tau} U(t)
\]

(1)

\[
\tau = D' \sigma_a^2 / c
\]

(2)

\[
D' = D_s \left(1 - D_s / D\right),
\]

(3)

where \(\sigma_a\) is the standard deviation in scattering angle \(a\) in each coordinate direction and \(c\) is the speed of light (Cordes & Lazio 2001). The unit step function, \(U(t)\), ensures that we only consider time \(t > 0\). Equation (1) is valid for a scattering screen that is infinite transverse to the line of sight. Such an infinite screen will conserve flux coming from the pulsar; every ray scattered out of the line of sight is compensated by another ray that hits the screen and is scattered back into the line of sight. For finite screens, the total flux of the source may not reach the observer depending on the physical extent of the screen and the observing frequency (see discussion in Section 3.3).

Anisotropic broadening functions. An anisotropic scattering screen would, in general, lead to the observation of elongated images of point sources behind such screens. These effects have, for example, been observed in the imaging of pulsar B0834+06 at 327 MHz (Brisken et al. 2010).

We consider an anisotropic scattering mechanism that can be modelled by an asymmetric Gaussian distribution in scattering angles. The temporal broadening function will now be dependent on two different characteristic scattering times, say \(\tau_x\) and \(\tau_y\), corresponding to the two dimensions that define the scattering screen, and has the form

\[
f_i = \frac{1}{\sqrt{\tau_x \tau_y}} e^{-x^2 / \frac{1}{2} \tau_x} I \left(\frac{1}{2} \left(\frac{1}{\tau_x} - \frac{1}{\tau_y}\right)\right),
\]

(4)

where \(I\) is the modified Bessel function of the first kind. The scattering times are defined as in equation (2) with \(\tau_{x,y} = D' \sigma_{a,x,y}^2 / c\) which reduces to the isotropic case for \(\tau_x = \tau_y\).

Numerical broadening functions. In Cordes & Lazio (2001), it was shown that the broadening function for a disc-shaped isotropic scattering screen centred on the line of sight can be created by amending equation (1) with a unit step function, such that the broadening function becomes \(f_i U(t_{\max} - t)\). This means that beyond a maximum observation angle \(\theta_{\max}\), no photons are scattered back into the observer’s line of sight. The maximum angle has a geometric path length correspondence to a maximum time delay \(t_{\max}\), as measured with respect to the direct line of sight.

We expand on these effects by means of a ray-tracing code that allows us to place a screen of any size between the source and the observer. We simulate a rectangular screen as a toy model: by specifying a screen’s width and height, its position along the line of sight \((D_s)\) and its 2D offset with respect to the line of sight, the scattering geometry is defined.

The temporal broadening function is constructed by calculating the probability that a given photon hitting the truncated screen will reach the detector. Binning these probabilities according to the associated arrival times for each photon allows us to build up the broadening function with a required time resolution.
2.2.3 Scattered profiles

Fig. 2 gives a schematic overview of two methods for producing scattered profiles. The method shown in the top panel represents the most intuitive way of creating folded scattered pulse profiles: a train of Gaussian pulses are convolved with a broadening function resulting in a train of broadened pulse shapes. A broadened pulse towards the end of the train (shaded) is representative of the average pulse profile. We refer to this method as the long train method.

An alternative method is shown in the bottom panel of Fig. 2. This method, which we simply call the train method, convolves a Gaussian pulse train of two pulses with a truncated temporal broadening function defined over a single pulse period (solid line). The second broadened pulse is then extracted as the averaged pulse profile (shaded). The outcome of the long train method (scaled by a constant factor and overplotted in dashed lines in the right-most frame) converges to the same average profile shape after a sufficient period of time. Fig. 3 shows the convergence of the two methods. The residuals between the last broadened pulse resulting from the train method and the long train method are shown to decrease as the pulse train length ($N$) in the latter is increased.

We note that the long train method requires care in the choice in the length of both the pulse train and the exponential, to ensure that the convolution converges to a constant set of pulses. Typically, the required pulse train length increases with an increase in scattering, thereby also increasing the computation time. The advantage of the train method is that it is both computationally less expensive and its accuracy is not dependent on the length of the broadening function or the pulse train (beyond 2 pulses). This motivates us to make use of the train method throughout this paper.

Finally, after generating the noiseless scattered profiles, Gaussian noise is added to simulate an observation with a chosen peak signal-to-noise ratio (SNR). Starting from an unscattered pulse with an off-pulse baseline of zero, all methods described above will, in the case of high levels of scattering (large $\tau$ compared to the period $P$), result in profiles with a raised off-pulse baseline. We use a Gaussian kernel to smooth the noisy scattered profiles. From the smoothed profile, we determine the raised baseline value to subtract. The impact of this on the fitting will be discussed in Section 3.1.

The frequency range over which we simulate data is chosen to create severe scattering imprints at the lowest considered frequency, where $\tau$ is a substantial fraction or even larger than the pulse period. In these cases, special care must be taken in the forward fitting process, as discussed in the next section. Although not presented in this paper, the simulations can be extended to include scattering screens with Kolmogorov (turbulent) scattering statistics.

2.3 Fitting techniques

Forward fitting techniques create models to fit the data by convolving an estimated intrinsic pulse shape with a parametrized...
broadening kernel that describes the response of the ISM (Ramachandran et al. 1997). In general, it is difficult to obtain a model for the intrinsic pulse shape as pulsars often exhibit frequency evolution. In this paper, we approximate the intrinsic pulse shape with a Gaussian function, with the caveats already stated.

An example of a forward fitting method is an exponentially modified Gaussian (EMG) method. In its most basic form, this gives the analytic expression for a single Gaussian pulse convolved with a one-sided exponential function. Such a simplistic implementation would not work well for large scattering times compared to the pulse period.

As mentioned earlier, inverse techniques perform a deconvolution of the broadened pulse shape through e.g. the CLEAN algorithm (Bhat et al. 2003). We focus on forward techniques. As shown in equation (2), the spectrum of a Gaussian isotropic scattering model has the form $\tau \propto v^{-4}$. We test the efficiency of fitting methods by determining how accurately they replicate this theoretical spectrum.

We use the lmfit package in PYTHON for our fits (Neville et al. 2014). We use our methods of generating noiseless scattered profiles, described above, as the fitting functions. The method fitting parameters are the underlying Gaussian intrinsic pulse components, i.e. the width ($\sigma$), mean ($\mu$) and amplitude ($A$), as well as the characteristic scattering time ($\tau$) and a DC offset, described in the next section.

The aim is to fit the simulated data for a range of frequencies to thereby obtain a spectrum of the characteristic scattering time, $\tau$. In doing so, we bear in mind that all observed scattered broadened profiles will typically result from integrating over some frequency band. The $\tau$ value obtained for fitting a monochromatic simulated scattered profile at frequency $f_m$ corresponds to the $\tau$ value associated with the central observing frequency $f_c$ and a bandwidth $\delta f$ by

$$f_c = 10^{\log_{10}(f_m + \delta f/2) + \log_{10}(f_m - \delta f/2)/2}$$

or inversely,

$$f_m = \frac{1}{2} \sqrt{\delta f^2 + 4 f_c^2}.$$  (6)

This equation provides the relevant frequency for a $\tau$ measurement, given any $f_c$ and $\delta f$, assuming a power-law spectrum. Therefore, computed $\tau$ values should be plotted against calculated $f_m$ values to obtain meaningful fits. We will revisit the impact of this in the next section. We note that in our methodology, we fit all frequency bands independently, whereas more constraints could be placed by simultaneously fitting for $\tau$ and the pulse parameters across a broad bandwidth, if that were available.

3 RESULTS

3.1 A simple end-to-end experiment

Before moving on to detailed results, we describe a simple end-to-end experiment. We simulate an intrinsic pulse shape scattered by an isotropic infinite scattering screen as described in Section 2.2. The frequency range over which we conduct the simulations is 60–140 MHz. Within this frequency range, the pulse profiles of the pulsar with $P = 1.0$ s and a duty cycle of 2.5 per cent, are notably scattered and in the severest of cases (lowest frequencies), the scattering tail wraps around.

Initially, we generate profiles monochromatically at 60 and 90 MHz with a peak SNR of 20. We fit the noisy profiles using the train method. We find that the determination of the off-pulse baseline of the noisy data has an important effect on the fit, which we investigate as follows. As mentioned, we use a Gaussian kernel to smooth the simulated data and from this smoothing function, estimate the off-pulse baseline, denoted in the following as $b$. The effect of the error in accurately determining the baseline, on the subsequently fitted $\tau$ value, is shown in Fig. 4. Here, we compare the fitted $\tau$ values to the outcome when the off-pulse baseline is known exactly. The plots show a percentage error in $\tau$, 100($\tau_{\text{known}} - \tau_{\text{fit}}$)/$\tau_{\text{known}$ plotted against the error in the determined baseline in units of the noise, ($b_{\text{known}} - b_{\text{fit}}$)/$\text{rms}_{\text{noise}}$. For 100 realizations of the noise at 60 and 90 MHz.

The error in the $\tau$ value increases systematically as the error in the baseline value grows, and is greater at the lower of the two frequencies. At 60 MHz, an error of 10–20 per cent in the $\tau$ value is already possible for baseline errors below $1 \sigma$ of the noise. As the frequency increases, the error in $\tau$ for a given error in the baseline drops significantly. We also note that at lower frequencies, the $\tau$ value is more likely to be underestimated, whereas at higher frequencies, $\tau$ is more likely to be overestimated.

This result prompted us to investigate an improved model. In Fig. 5, we show the impact of adding a DC offset to the parameter set for which the model fits. In the train + DC model, we therefore no longer estimate the off-pulse baseline in order to subtract it
Figure 5. Histograms showing the obtained spreads in $\tau$ values at three sample frequencies, using different methods. The red line indicates the true value of $\tau$. The best check histogram shows the distribution obtained from fits where the noiseless minimum of the profile is known exactly, and is therefore an idealized reference case. The train and train + DC methods show the spreads in $\tau$ using the train method with or without fitting for a DC offset. In the case where the method does not fit for a DC offset, a Gaussian smoothing kernel is used to estimate and subtract the raised off-pulse baselines.

Figure 6. The confidence intervals (1–3$\sigma$) associated with the obtained fits for the vertical offset parameter, DC, and characteristic scattering time, $\tau$. There are deviations from the theoretical characteristic $\tau$ spectrum observed as the bandwidth of the simulated data increases. The obtained $\tau$ associated with a central frequency $f_c$ and bandwidth $\delta f$ relate to the simulated result produced by a monochromatic frequency input $f_m$ through equation (5). Fig. 7 shows the impact on the $\tau$ spectrum fit when the central frequency values $f_c$ are used instead of $f_m$. Simulating the data monochromatically at 20 MHz intervals from 60 to 140 MHz leads to a measured spectral index of $\alpha = 4.00 \pm 0.04$, reproducing the theoretical spectrum well. In contrast, a bandwidth of 20 MHz at 20 MHz intervals gives $\alpha = 3.93 \pm 0.04$, when plotted against $f_c$. An even larger bandwidth of 40 MHz gives $\alpha = 3.70 \pm 0.04$. To produce this figure, data were fitted with the train method and the baseline of the simulated data subtracted before adding noise – to eliminate the effects of baseline errors as mentioned above.

3.2 Results for infinite screens

We now move on to infinite scattering screen setups. We first present the obtained $\tau$ spectra, followed by the flux spectra. In each case, we consider the two scattering mechanisms discussed in Section 2.2.2, which serve as examples of either an isotropic or anisotropic scattering process.
Figure 8. Best fits as produced by the train + DC method for simulated scattered pulse profiles of a $P = 1.0$ s pulsar. Fits without the DC-offset parameter (i.e. the train method in which the off-pulse baseline is determined using a smoothing kernel) look similar to the DC fits by eye, but produce skewed values for $\tau$, as was shown in Fig. 5. The impacts of these offsets on the power-law fit to the $\tau$ spectra can be seen in Table 1. Profiles are plotted normalized to the area under the pulse – such that at the lowest frequency (viz. 60 MHz) in the absence of a non-zero off-pulse baseline, the area under the pulse is equal to one.

3.2.1 $\tau$ spectra

3.2.1.1 Isotropic scattering $\tau$-spectra. Isotropically scatter broadened profiles of the 1.0 s simulated pulsar, described in Section 2.2.1, are shown in Fig. 8. The simulated data, as seen in yellow, again have peak SNR = 20. The scattering setup is as described in Section 2.2.2. A range of observing frequencies is chosen such that the scattering tail wraps around at the low end and a typical exponential scattering tail is observed at the high end.

The profiles are fitted with the train + DC method as described in the previous section. Fits without the DC-offset parameter (i.e. the train method in which the off-pulse baseline is determined using a smoothing kernel) look very similar to the DC fits by eye, but produce skewed values for $\tau$ as was shown in Fig. 5. The impacts of these offsets on the power-law fit to the $\tau$ spectra can be seen in Table 1.

As described in Section 2.2.3, the train and the long-train methods, with added DC offsets, can perform equally well for all the simulations. A long train method with a too short pulse train will lead to an underestimation of $\tau$.

In the case of severely scattered profiles, it is vital to use an accurate fitting method as the outcome becomes increasingly sensitive to the model parameters. The train + DC method presents the simplest (and computationally fastest) way to obtain accurate $\tau$ values.

The $\tau$ spectra corresponding to Fig. 8 are shown over two frequency ranges in Fig. 9. Similarly the fitted profiles and $\tau$-spectra for the 20 ms pulsar are shown in Figs 10 and 11. For both pulsars, as the frequency increases, the obtained spectral index value more accurately approaches the simulated index of $\alpha = 4$. Here we also include a basic implementation of the EMG model not accounting for the pulse period $P$. This method underestimates $\tau$ significantly at high levels of scattering. At higher frequencies where the scattering tail no longer wraps around, the choice of method becomes less important.

We generate 1000 data sets, of simulated scattered profiles at intervals of 5 MHz within our frequency ranges, per pulsar and run the fits for $\tau$ and $\alpha$. The outcomes of this process are summarized in Table 1.

For each realization of the experiment, the train + DC method performs best, and, in the case of the millisecond pulsar (MSP; for which the scattering is more severe overall), has a lower standard deviation in its estimation of $\alpha$ than the standard train method. The average spectral index is calculated over two sets of frequencies for both the slow and MSP. At higher frequencies (lower scattering), $\alpha$ approaches the simulated value of 4 more accurately. This can for example be seen in the tabulated error values, which represent the error of the mean spectral index with respect to the true value of 4 expressed as a percentage.

In Fig. 5, we showed the spread of $\tau$ values obtained by the train methods and by a best check reference case. As noted, even the
Table 1. The mean spectral indices and standard deviations as obtained from 1000 τ and subsequent spectra fits. Two pulsars are investigated: a slow pulsar and a MSP. For each pulsar, the spectral index is determined over two sets of frequencies. The fits and their associated standard deviations improve with a decrease (increase) in scattering (frequency). The different methods are as described in the text.

| Frequency Range: 60 – 120 MHz | EMG | Train | Train + DC | Train + DC asymm errors |
|--------------------------------|-----|-------|------------|------------------------|
| Mean Spectral Index            | 2.90| 3.80  | 4.00       | 4.01                   |
| Standard deviation             | 0.05| 0.09  | 0.09       | 0.09                   |
| % Error                        | −27.6%| −4.9%| −0.1%    | 0.2%                   |
| Fast pulsar: 20 ms, 10% duty cycle SNR: 20 |
| Mean Spectral Index            | 3.78| 3.89  | 4.00       | 4.00                   |
| Standard deviation             | 0.06| 0.06  | 0.06       | 0.03                   |
| % Error                        | −5.4%| −2.8%| 0.0%      | 0.0%                   |
| Frequency Range: 90 – 150 MHz |
| Mean Spectral Index            | 1.11| 3.71  | 3.87       | 4.01                   |
| Standard deviation             | 0.06| 0.42  | 0.32       | 0.17                   |
| % Error                        | −72.2%| −7.2%| −3.1%    | 0.2%                   |
| Frequency Range: 200 – 260 MHz |
| Mean Spectral Index            | 2.38| 3.85  | 3.98       | 3.99                   |
| Standard deviation             | 0.09| 0.18  | 0.14       | 0.07                   |
| % Error                        | −40.4%| −3.7%| −0.5%    | −0.3%                  |

The train + DC model shows a skewed distribution of τ especially at low frequencies. To take this skewness into account, we update the power-law fit as follows. From a distribution of τ values (obtained from several noise realisations at a given SNR level, and subsequent profile fits with the train + DC method), we subtract 16 per cent of the values from both the lower and upper end of the distribution to be left with the 1σ (68 per cent) central distribution. We now use the values within this range along with their corresponding frequency values to update the τ spectrum, as shown in Fig. 12. As can be seen from Table 1, under the column heading Train + DC, asymm errors, this process mainly provides a small improvement at extreme levels of scattering, such as for the MSP at 150–210 MHz. For the most part of this article, we will therefore not show these improvements.

As is expected, the success of the τ spectra fits and the associated uncertainties are also impacted by the peak SNR value, and will improve as the SNR increases. In Table 2, we show the achieved spectral index values after repeating the experiment above at SNR values of 10 and 50. At a SNR value of 50, the theoretical spectrum is best reproduced and the standard deviation in the spread of individual measurements has decreased significantly.

3.2.1.2 Anisotropic scattering τ-spectra. In the case of the anisotropic scatterer discussed here, the broadened profile is dependent on two characteristic scattering times, each with a frequency dependence τ_x, y ∝ ν^-4, as described in Section 2.2.2. In general, a more complex frequency dependence could be at play, since different observing frequencies will interact with anisotropic scattering ellipses of varying sizes. Sampling different regions of the scattering screen, in this way, can lead to additional frequency dependences not included in our model.

We show how an anisotropic scattering process will create apparent deviations from the expected τ spectrum when fitted with an isotropic scattering model. The scattering strengths are chosen to be σ_ax = 3 mas and σ_ay = 1 mas at 1 GHz. That is to say the screen scatters three times more weakly in the chosen y-dimension than in the x-dimension. Setting σ_ay = 3 mas will reproduce the isotropic scatterer considered in Section 3.2.1.1. The scattering geometries are as before and the pulsar is the same 1 s Gaussian profile from the previous section.

The simulated data are fitted with the train + DC method. The resultant profile fits for a selection of frequencies are shown in Fig. 13. The corresponding τ spectrum is shown in Fig. 14. As before, these are chosen to represent the mean profiles and spectrum after repeating the experiment with Gaussian noise of peak SNR = 20 multiple times. The mean and standard deviation of the spectral index as calculated after 1000 repetitions are given in Table 3.
Figure 10. Similar to Fig. 8, but for a pulsar with period $P = 20$ ms and a duty cycle of 10 per cent.

Figure 11. The $\tau$ spectra of the modelled MSP. The uncertainty in the spectral index obtained by the train + DC method grows as the frequency range is lowered. At these lower frequencies, the EMG method deviates most severely from the theoretical spectrum, and continues to underestimate the value of $\alpha$ even for the higher frequency range. The frequency at which the choice of method matters is much higher for the simulated MSP than for the slow pulsar in Fig. 9.

Table 2. Spectral indices similarly obtained as for Table 1. The simulated data have SNR values of 10 and 50, and show the expected improvements in fits as the SNR values increase. The tabulated values represent a mean $\alpha$ and associated standard deviation after 1000 executions of the experiment.

In Fig. 14, the input values of $\tau_1$ and $\tau_2$ (green solid and dashed lines) are shown as a reference. Both have a frequency dependence with spectral index $\alpha = 4$. The geometric mean,

$$\tau_{geo} = \sqrt{\tau_1 \tau_2},$$

which would reduce the anisotropic scattering to isotropic scattering for $\tau_1 = \tau_2$, is computed from the input $\tau_1$ and $\tau_2$ and shown in yellow (stars). The single $\tau$ fit, shown in red, provides an estimate of the $\tau_{geo}$ value at mid-frequencies. At low frequencies, however,
Figure 13. Profile fits as obtained by applying the isotropic train + DC method to anisotropically simulated pulse profiles of a $P = 1.0$ s pulsar. The anisotropy is created by choosing a thin screen that scatters three times more strongly in one dimension than in the other. At higher frequencies, the isotropic fit to anisotropic data leads to a rough estimation of the geometric mean scattering time ($\tau_{\text{geo}}$). At low frequencies, $\tau_{\text{geo}}$ is significantly underestimated. Profiles are plotted normalized to the area under the pulse, such that at the lowest frequency (viz. 60 MHz) in the absence of a non-zero off-pulse baseline, the area under the pulse is equal to one.

Table 3. The mean spectral indices as obtained by fits to profiles produced through an anisotropic scattering process. The tabled spectral indices represent the average index as obtained from 1000 fitting procedures.

| Anisotropic Scattering |
|-------------------------|
| Slow pulsar: 1.0 sec, 2.5% duty cycle |
| SNR: 20 |
| Frequency Range: |
| 60 – 120 MHz |
| Train + DC |
| Mean Spectral Index | 2.95 |
| Standard deviation | 0.05 |
| % Error | -26.3% |

Figure 14. The $\tau$ spectrum for an anisotropic scattering event fitted with a single $\tau$ power law is shown in red (thick dashed line, circles). The green lines (solid and dotted without markers) show the simulated values for $\tau_1$ and $\tau_2$. The geometric mean for each of these values, $\tau_{\text{geo}} = \sqrt{\tau_1 \tau_2}$ is shown in yellow (stars). The $\tau_1$, $\tau_2$ and $\tau_{\text{geo}}$ values have a modelled frequency dependence $\tau \propto \nu^{-4}$. At low frequencies, the fits deviate most significantly from the $\tau_{\text{geo}}$ values.

These results show that scatter broadened profiles caused by anisotropic scattering mechanisms require more sophisticated fitting models that can account for the degree of anisotropy. Such a model would make explicit use of equation (4) in its description. For individual pulsars, screen anisotropy might be revealed through elongated images (Brisken et al. 2010) or features of the secondary spectra, which will be instructive in determining the choice of fitting function, and in improving the fit.
3.2.2 Flux-spectra

The mean flux ($S_m$) of a pulse profile is calculated by integrating the intensity of the profile averaged over the pulse period. This is illustrated schematically on the right-hand side of Fig. 2, where the mean observable flux is shaded in blue. Scatter broadening that causes the power of pulses to be smeared into succeeding pulses will lower the observable flux, as intrinsic flux is lost to the raised baseline. Since this is an artefact of the temporal domain, we can, in principle, circumvent the flux loss by imaging pulsars directly (e.g. Dembska et al. 2015).

We compute the flux for the setups in Section 3.2.1 using the 1 s pulsar, and show that a turnover in the flux spectrum occurs at low frequencies (high scattering). The location of this turnover is dependent on scattering mechanism and geometry.

The mean flux values quoted here are calculated from the noiseless scatter broadened pulse shapes. It is normalized such that in the absence of the baseline being raised, the lowest considered frequency (viz. 40 MHz) will have a mean flux of unity.

Fig. 15 show the turnovers in spectra for both the isotropic and anisotropic scattering setup. The turnover in the flux spectrum for the anisotropic scattering case occurs at a lower frequency and causes a smaller loss in flux since the mean scattering is weaker. The turnover also has an asymmetric shape.

Fits to the high-frequency end of these spectra approximate the intrinsic spectral index of $\gamma$ = 1.6. For low frequencies, however, the spectral turn-around leads to a positive slope, such that the spectral turnover appears. Simultaneously, as $D$ increases, the frequency window within which spectral turnover takes place for all plotted ratios of $D_s/D$ (indicated with dashed lines) grows slowly. This result shows that for distances within our Galaxy and a scattering strength of $\sigma_s = 3$ mas at 1 GHz, as taken from Cordes & Lazio (2001), turnovers in mean flux spectra will occur below 100 MHz. Spectral turnovers at higher frequencies are either due to different scattering mechanisms (deviations from isotropically Gaussian scattering) or different scattering strengths than the typical value used here. However, since a wide range of scattering strengths are expected, it is unclear whether or not a coherent picture is expected throughout the population of pulsars with measured spectra.

At very large distances, $D > 200$ kpc, (while still using the above scattering strength) the simulated pulse is completely scattered out and a turnover in the flux spectrum no longer detectable in this way.

3.3 Effects of screen geometry

In Section 2.2.2, we discussed computing the temporal broadening functions ($f_s$) of truncated screens by means of a ray-tracing code.
We expect these finite screens to cause a loss in the observable flux of the pulsar. Furthermore, we investigate how the shape of an average pulse profile, and consequently the $\tau$ spectrum, is impacted by the presence of such a screen.

By defining the screen in terms of its height and width, we can create both square screens that will have a comparable impact to a circular screen or rectangular (elongated) screens. In the theoretical limit where screens become infinitely more long than wide (i.e. 1D circular screen or rectangular (elongated) screens. In the theoretical limit where screens become infinitely more long than wide (i.e. 1D circular screen or rectangular (elongated) screens, the broadening function will tend to,

$$f_1 = \frac{e^{-t/\tau}}{(\sqrt{\pi} \tau)} U(t).$$

We stick to screens where the ratio of its dimensions has not yet reached this theoretical limit.

### 3.3.1 Profile effects

The broadening function for a source at a distance of $D = 3$ kpc with a finite isotropic scattering screen halfway between the source and the observer is presented in Fig. 17. The modelled screen has size 400 au $\times$ 600 au, with the pulsar offset to the middle of the screen by 100 au horizontally and 50 au vertically. The scattering strength is as before, $\sigma_a = 3$ mas at 1GHz.

The dips in the broadening function are associated with the four edges of the screen and show the associated loss in flux at different frequencies. In this instance, we see dips at 0.02, 0.10, 0.15 and 0.20 s associated with distances to the edge of 100, 250, 300 and 350 au.

The numerical broadening functions are convolved with the same $P = 1.0$ s pulsar as used for the infinite scattering screens in Section 3.2. A selection of the resulting broadened profiles are shown on the left-hand side of Fig. 18.

The broadening effects can be split into three regimes. At high frequencies (above 210 MHz, not shown), the effect of scattering is negligible and the observed pulse resembles the intrinsic Gaussian pulse. At mid-frequencies (here 120–210 MHz), the pulse shapes have the typical exponential tail as would be associated with an infinite scattering screen. At low frequencies (below 120 MHz), the finite nature of the scatterer becomes visible in the form of bumps associated with the loss in flux at various path lengths (or times). These pulse shapes have a much steeper trailing edge than normally associated with a scatter broadened pulse shape.

As an example of these effects on a more complicated profile, we use a high-frequency profile of B1237+25 as a multicomponent template and consider how this template would change with frequency when it is scattered by a truncated screen. In our simulations, we use the pulsar’s period ($P = 1.38$ s) and distance ($D = 0.85$ kpc), while acknowledging that the pulsar itself, being a high galactic latitude source, is not an ideal candidate for scattering studies. We consider the impact of an isotropic screen, size 200 au $\times$ 200 au, placed mid-way along the line of sight. The scattering strength is as before. The resulting broadened pulse shapes at several frequencies are shown on the right-hand side of Fig. 18.

Note how the truncation effect of the screen can cause the height of the two profile peaks to change relative to each other. Between 110 and 130 MHz, the left peak is lower than the right one, at 140 MHz, the peaks are roughly equal and beyond 140 MHz, the left peak is the higher peak. The change in features in these profiles could easily be misinterpreted as features of the intrinsic pulse shape.

The above scattering setups are theoretical constructs, however, by investigating the impact of truncated screens on both the $\tau$ and flux spectra, we can build a more complete picture of all the possible sources of deviations to the otherwise theoretically predicted spectral trends.
Scattering imprints on pulsar observations

Figure 19. Profiles as generated by the truncated broadening function of Fig. 17 with added Gaussian noise. The profiles plotted here are normalized relative to the peak of the profile at 40 MHz (which has suffered the greatest flux loss). The input $\tau$ values at different frequencies are given, along with the values obtained from the fits with the train method. In all instances, the $\tau$ values are underestimated by more than 45 per cent.

3.3.2 Effects on $\tau$ spectra

To estimate the $\tau$ spectra, we added Gaussian noise to the simulated profiles, after which we fitted them with the train method of Section 2.3. (Since the truncation prevents wrap-around profiles here, we do not have to include a DC-offset parameter.) The SNR value is modelled to scale with the mean flux of the profile, and is therefore notably smaller at low frequencies for which the flux loss is greater. (We discuss the flux spectrum in more detail in the next section.) The resulting profile fits for a selection of frequencies are shown in Fig. 19.

The corresponding $\tau$ spectrum is shown in Fig. 20. From the $\tau$ spectrum, it is clear that a model based on the assumption of an infinite screen, greatly underestimates the true value of $\tau$. A weighted fit to all the obtained $\tau$-values produces a spectral index of $\alpha = 3.17$. If we identify a single frequency at which the observations appear to become sensitive to the edges of the screen, we can fit the spectrum with two power laws instead. Using 100 MHz as this breakpoint, we find that $\alpha = 2.85$ above this frequency. Below 100 MHz, we see a flattening of the $\tau$ spectrum, such that the obtained spectral index here is $\alpha = 0.34$.

Figure 20. The $\tau$ spectrum associated with a truncated screen. At low frequencies (high scattering), the $\tau$ values as obtained by the train method (red circles, with error bars smaller than the marker) greatly underestimates the input scattering (yellow triangles). This is due to the fact that at large wavelengths, the scattering process will become sensitive to the edges of the screen leading to imprints in the temporally broadened pulse profiles for which our method does not account. The short black (dashed) and blue (dot–dashed) lines show a broken power-law fit with spectral indices as indicated in the legend, and the long red (dashed) line is the weighted power-law fit of all the obtained $\tau$ values.
Improvements to such τ value is centred on the true simulated value also contains the flux spectrum for a truncated anisotropic τ 200 au). That it is important to attribute the measured γ increase towards lower 1 at the lowest observed frequency (provided that no flux is included an example (as a grey line) of a τ estimates. In the case of the 1.0 s pulsar at 60–120 MHz, γ 16 becomes 30–70 per cent for such an implementation, deviations from the theoretically expected values at low frequencies can be as large as 30–70 per cent for the pulsars and scattering scenarios we simulate. The failure of this model is made explicit to encourage the revision of fits in the literature where wrap-around scattering tails may not have been modelled, e.g. fig. 1 in Lewandowski et al. (2013).

We show in Fig. 7 that it is important to attribute the measured τ values to the associated monochromatic frequency (f_m, computed from the central frequency and bandwidth of the observation as in equation 6) to obtain accurate estimates of the spectral index (α) from a power-law fit.

We investigate the scatter broadening of pulsars due to anisotropic scattering, which we model as an asymmetric Gaussian distribution, and show that fitting these profiles with isotropic models leads to systematic errors in the measurement of τ. Assuming the scattering strength in one dimension to be weaker than what we have considered as typical in our simulations, such fits will lead to apparent τ spectra with α < 4 (in our example α = 2.95). Improvements to such fits will rely on using more appropriate broadening functions, such as equation (4) or (8), in the fitting model. We are currently investigating the use of anisotropic fitting techniques to analyse scattered profiles from LOFAR observations. Given the recent evidence for highly anisotropic scattering, our simulations show that the measured α (and goodness of fit for τ) could serve as a diagnostic in the time domain for the type of scattering screen.

From the simulated truncated screens in Section 3.3, we show that an isotropic thin screen with dimensions 400 au × 600 au midway between the Earth and a pulsar at 3 kpc, changes the shape of keeping the scattering strength, could be by means of ever smaller scattering clouds. Fig. 21 includes an example (as a grey line) of a smaller screen with a higher turnover frequency.

4 DISCUSSION

Reflecting on the results in the Section 3, we make the following statements.

In Section 3.1, we show that errors related to estimating the off-pulse baseline can be circumvented by treating the baseline level (DC offset) as a parameter for which our methods fit. This ensures that the estimated τ value is centred on the true simulated value for τ.

Table 1 summarizes the impact of using different methods to estimate the spectral indices. Errors in τ increase towards lower frequencies (or higher scattering) for all the fitting methods. Given a pulse shape and a duty cycle, τ in MSP will become comparable to the pulse period at higher frequencies than for normal pulsars.

Larger errors in τ lead to larger errors in α. Our simulations show that at high levels of scattering, there is an underlying asymmetry in the measurement uncertainties of τ. Correctly accounting for the asymmetry results in a more accurate measurement of α. The impact of this correction is seen for the MSP case, where the mean estimate of α changes from 3.87 to 4.01 over the low-frequency range, while the standard deviation in the estimate simultaneously drops from 0.32 to 0.17. Having introduced this correction, the α values for all our simulated data have a mean value within 0.3 per cent of the true value, and a standard deviation below 0.2.

Increasing the SNR also leads to a decrease in the standard deviation of α estimates. In the case of the 1.0 s pulsar at 60–120 MHz, an increase of SNR from 10 to 20 to 50 decreases the standard deviation by a factor of 2 and 2.25, as shown in Table 2.

The EMG method in this paper portrays the most basic model, where the wrap-around of a scattering tail is not modelled. For such an implementation, deviations from the theoretically expected values at low frequencies can be as large as 30–70 per cent for the pulsars and scattering scenarios we simulate. The failure of this model is made explicit to encourage the revision of fits in the literature where wrap-around scattering tails may not have been modelled, e.g. fig. 1 in Lewandowski et al. (2013).

We show in Fig. 7 that it is important to attribute the measured τ values to the associated monochromatic frequency (f_m, computed from the central frequency and bandwidth of the observation as in equation 6) to obtain accurate estimates of the spectral index (α) from a power-law fit.
the pulse profile with frequency and leads to a flattened $\tau$ spectrum below 200 MHz. In Fig. 20, the spectral index was estimated to be $\alpha = 0.34$ below 100 MHz. Adjustments to the modelled screen size and location will impact the shape and features of the $\tau$ spectra, resulting in possible breaks.

For a complex pulse profile, the frequency evolution due to this kind of scattering may change the component amplitude ratio. We use the profile of B1237+25 to demonstrate this effect, although this pulsar is not observed to exhibit such behaviour. The effects caused by scattering can, in principle, be disentangled from the intrinsic emission, as it is unlikely that both will have the same dependence on frequency.

Fig. 22 provides a cartoon style summary of deviations from the theoretical $\tau$ spectrum we use (a power-law spectrum with spectral index 4) and the underlying causes. Subscripts $I$ and $A$ refer to isotropic and anisotropic scattering mechanisms, and $\infty$ or $tr$ to either infinite or truncated thin screens.

For the range of scattering mechanisms and geometries discussed above, we also investigated the impact on pulsar flux spectra. Observationally, pulsar flux spectra have steep negative spectral indices, typically within the range $\gamma = -1.8$ to $-1.4$ (Lorimer et al. 1995; Maron et al. 2000; Bates, Lorimer & Verbiest 2013). Spectral turnovers at low frequencies are a known feature of the population of slow pulsars (e.g. Lorimer et al. 1995). Thermal absorption by dense surrounding environments (e.g. Rajwade, Lorimer & Anderson 2016) may be partially responsible for spectral turnovers, which may also be intrinsic to the radio emission of the pulsar. Some evidence suggests that the efficiency of the radio emission shuts off at lower frequencies (Malofeev, Malov & Shchegoleva 2000). Alternatively, under the assumption that the lower frequencies are emitted higher above the neutron star surface than the higher frequencies (Radhakrishnan & Cooke 1969), the intersection of the line of sight with the radio-emitting regions in the magnetosphere will also be frequency dependent, with the possibility of the line of sight gradually missing the emitting regions altogether at low frequencies. A small group of pulsars show turnovers at much higher frequencies. These are known as the gigahertz peaked spectrum (GPS) pulsars. The GPS turnovers are also attributed to either the radio emission process becoming inefficient above a threshold frequency, or dense ISM environments (Kijak, Gupta & Krzeszowski 2007; Kijak et al. 2011).

Here we concentrate on the effects of scattering on the observed flux, which should occur in addition to the above mechanisms. In our simulations, we show that high levels of scattering by infinite screens result in a measurable integrated flux loss, due to the rise in the off-pulse baseline. In the case of slow pulsars, this will lead to spectral turnovers typically below 100 MHz, as shown in Fig. 16. The turnover frequency, due to this effect in MSP flux spectra, would be higher than for slow pulsars. Older studies of MSP flux spectra have mostly not found turnovers (Kuzmin & Losovsky 2000; Kuzmin 2002); however, more recent work by Kuniyoshi et al. (2015) have observed turnovers in MSP flux spectra, albeit in imaging observations. This could point to anomalous scattering by e.g. finite scattering clouds. The lack of observed spectral turnovers from time-domain studies suggest that most MSP profiles are not substantially scattered in current low-frequency observations.

Fig. 15 shows the difference in turnover frequencies for an infinite isotropic and anisotropic scattering screen. In our example of an anisotropic screen, the spectral turnover occurs at $\sim$50 MHz. The turnover frequency for the corresponding isotropic case was $\sim$90 MHz, i.e. higher by a factor of $\sqrt{\sigma_{\text{iso}}/\sigma_{\text{i}}}$ (where subscripts $a$ and $i$ refer to anisotropic and isotropic, respectively, and $(x, y)$ are the two dimensions describing the thin screen, with $\sigma_{a} = \sigma_{i} = 3$ and $\sigma_{a} = 1$ mas at 1GHz).

Similar spectral turnovers can be caused by the loss of flux associated with truncated scattering screens. In Fig. 21, we show that a finite screen of size 400 au $\times$ 600 au at the same location as the infinite screen, would lead to turnover frequencies in approximately the same range, with the isotropic flux spectrum peaking at 125 MHz and the anisotropic spectrum at 90 MHz. Keeping the scattering strength constant, the turnover frequency can be pushed up to 230 MHz for a screen of size 100 au $\times$ 200 au, centred on the line of sight. In this paper, we have concentrated on screens with infinitely small inner scales, i.e. the scales of structure in the screen that result in the largest scattering angles. A larger inner scale, as discussed in Rickett et al. (2009), would result in a maximum scattering angle, similar to the case of a screen of finite size, and hence similar spectra for $\tau$. The crucial difference between the two is that the existence of an inner scale does not reduce the total observed flux, as is the case for a finite screen.

To reproduce GPS spectra with our modelling would require a fine tuning of the truncated screen size and or the simulated scattering strength. Studying the correlated changes in flux spectra, pulsar profiles and their associated $\tau$-spectra with frequency can reveal the complex interplay between the phenomena that set the observed pulse profile and its dependence on frequency, and should lead to a better understanding of the effects of each.

5 CONCLUSIONS

We have showed that our train + DC method is an efficient forward fitting technique for scattered profiles. Our investigations show flat $\tau$ spectra (than what is theoretically expected) when using an isotropic model to fit profiles shaped by an anisotropic scattering mechanism. Flatter $\tau$ spectra can also be the result of low-frequency profile shape changes caused by a truncated scattering screen. Over a large enough sampled frequency range, such a $\tau$ spectrum will exhibit a break. Analysing the impact of scattering on flux spectra has revealed turnovers at low frequencies due to baseline effects. Similarly, truncated screens can cause turnovers in flux spectra.

![Figure 22: A cartoon summary of the deviations from the theoretical $\tau$ spectrum and their underlying causes. Subscripts $I$ and $A$ refer to isotropic and anisotropic scattering mechanisms, and $\infty$ or $tr$ to either infinite or truncated thin screens. The anisotropic scatterer shown here scatters more weakly in one dimension than what is considered a typical isotropic scattering strength.](https://academic.oup.com/mnras/article-abstract/462/3/2587/2589386)
These turnovers are a function of the screen size and its scattering strength.

We anticipate new data from instruments such as LOFAR and the MWA in the frequency range studied here. This paper will serve as a framework to compare real data to the simulated results obtained here. Our analysis of scattering imprints is by no means complete, as we have only investigated simple profile shapes and Gaussian scattering mechanisms. However, our results indicate where attention should be paid in terms of creating accurate fitting methods as well as describing possible sources of deviations to theoretically predicted results.

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