INTEGRABLE 1D TODA CELLULAR AUTOMATA

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Abstract. First, we recall the algebro-geometric method of construction of finite field valued solutions
\[ Z^3 \ni (n_1, n_2, n_3) \mapsto \tau(n_1, n_2, n_3) \in \mathbb{F}_q, \]
of the discrete KP equation (\( T_i \) — shift operator in \( n_i \) variable)
\[ (T_1 \tau) \cdot (T_2 T_3 \tau) - (T_2 \tau) \cdot (T_3 T_1 \tau) + (T_3 \tau) \cdot (T_1 T_2 \tau) = 0, \]
and next we perform a reduction of dKP to the discrete 1D Toda equation
\[ (T_1 T_3 \tau) \cdot (T_1 T_3^{-1} \tau) - (T_1^{-1} \tau) \cdot (T_1 \tau) + \tau \cdot \tau = 0. \]
This gives a method of construction of solutions of the discrete 1D Toda equation taking values in the finite field \( \mathbb{F}_q \).

1. Introduction

Discrete integrable systems focus much attention for at least two reasons. They inherit developments of the theory of continuous integrable models on the one hand, and give rise to new ideas on the other. Recently this ”twofold” nature was exploited for the construction of integrable cellular automata.

The approach is based on transferring the standard algebro-geometric method of construction of solutions [9, 1] from the complex field \( \mathbb{C} \) to a finite field case. In spite of apparent differences between complex and finite fields (see [11]) the theory of algebraic functions over finite field [7, 12] provides a powerful framework and makes this transfer possible. As a consequence of algebraic approach it appears that relevant theorems (as the Riemann-Roch theorem) and formulas are ”quite the same” as in the complex field case but have a different meaning. For instance, the concept of a continuum limits (in expansions etc.) is not clear as far as finite fields are regarded.

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For a given difference equation, its finite field valued solution defines a cellular automaton. So far there have been constructed multisoliton solutions for the discrete 2D Toda lattice (the Hirota’s bilinear difference equation) [6], the discrete Kadomtsev-Petviashvili (dKP) and discrete Korteweg-de Vries (dKdV) equations [3] and also algebro-geometric solutions for dKP [4] and dKdV [2]. In this article we extend previous results for the discrete 1D Toda equation (and ipso facto for a certain discrete version of the sine-Gordon equation, see remark at the end of section 3.1).

The paper consists of three main parts. In Section 2 we recall the finite field version of the algebro-geometric construction of solutions of the discrete KP equation. In Section 3 the algebro-geometric reduction to the discrete 1 dimensional Toda equation is performed and illustrated by examples in Section 4.

2. An algebro-geometric approach to dKP equation

2.1. General construction for the dKP equation. For the general construction [3] we need an algebraic projective curve $\mathcal{C}/\mathbb{K}$ (or simply $\mathcal{C}$), absolutely irreducible, nonsingular, of genus $g$, defined over the finite field $\mathbb{K} = \mathbb{F}_q$ with $q$ elements. $\mathcal{C}(\mathbb{K})$ denotes the set of $\mathbb{K}$-rational points of the curve. $\overline{\mathbb{K}}$ denotes the algebraic closure of $\mathbb{K}$, i.e., $\overline{\mathbb{K}} = \bigcup_{\ell=1}^{\infty} \mathbb{F}_{q^\ell}$, and $\mathcal{C}(\overline{\mathbb{K}})$ denotes the corresponding infinite set of $\overline{\mathbb{K}}$-rational points of the curve. Denote by $\text{Div}(\mathcal{C})$ the abelian group of the divisors on the curve $\mathcal{C}$. The action of the Galois group $G(\mathbb{K}/\mathbb{F}_q)$ (of automorphisms of $\mathbb{K}$ which are identity on $\mathbb{F}_q$) extends naturally to action on $\mathcal{C}(\overline{\mathbb{K}})$ and $\text{Div}(\mathcal{C})$. A field of $\mathbb{K}$-rational functions on the curve $\mathcal{C}$ is denoted by $\mathbb{K}(\mathcal{C})$. A vector space $L(D)$ is defined as $\{f \in \mathbb{K}(\mathcal{C}) \mid (f) > -D\}$, where $D \in \text{Div}(\mathcal{C})$ and $(f) = \sum_{P \in \mathcal{C}} \text{ord}_P(f) \cdot P$ is the divisor of the function $f \in \mathbb{K}(\mathcal{C})$.

On the curve $\mathcal{C}$ we choose:

1. four points $A_0, A_i \in \mathcal{C}(\mathbb{K})$, $i = 1, 2, 3$,
2. effective $\mathbb{K}$-rational divisor of order $g$, i.e., $g$ points $B_\gamma \in \mathcal{C}(\overline{\mathbb{K}})$, $\gamma = 1, \ldots, g$, which satisfies the following $\mathbb{K}$-rationality condition
   $$\forall \sigma \in G(\overline{\mathbb{K}}/\mathbb{K}), \quad \sigma(B_\gamma) = B_{\gamma'},$$
3. effective $\mathbb{K}$-rational divisor of order $N$, i.e. $N$ points $C_\alpha \in \mathcal{C}(\overline{\mathbb{K}})$, $\alpha = 1, \ldots, N$, that satisfies the $\mathbb{K}$-rationality conditions
   $$\forall \sigma \in G(\overline{\mathbb{K}}/\mathbb{K}), \quad \sigma(C_\alpha) = C_{\alpha'}.$$
(4) $N$ pairs of points $D_\beta, E_\beta \in C(\mathbb{K})$, $\beta = 1, \ldots, N$, that satisfy $\mathbb{K}$-rationality conditions

$$\forall \sigma \in G(\mathbb{K}/\mathbb{K}) : \sigma(\{D_\beta, E_\beta\}) = \{D_{\beta'}, E_{\beta'}\}.$$ 

We assume that all the points used are distinct and in general position. In particular, the divisors $\sum_{\gamma=1}^{g} B_\gamma$, $\sum_{\alpha=1}^{N} C_\alpha$, and $D(n_1, n_2, n_3)$, defined below, are non-special.

**Definition.** Fix $\mathbb{K}$-rational local parameter $t_0$ at $A_0$. For any integers $n_1, n_2, n_3 \in \mathbb{Z}$ let divisor $D(n_1, n_2, n_3)$ be of the form

$$D(n_1, n_2, n_3) = \sum_{i=1}^{3} n_i(A_0 - A_i) + \sum_{\gamma=1}^{g} B_\gamma + \sum_{\alpha=1}^{N} C_\alpha.$$ 

The function $\psi(n_1, n_2, n_3)$ (called a wave function) is a rational function on the curve $C$ with the following properties

1. the divisor of the function satisfies $(\psi) > -D$, i.e. $\psi \in L(D)$,

2. the function $\psi$ satisfies $N$ constraints

$$\psi(n_1, n_2, n_3)(D_\beta) = \psi(n_1, n_2, n_3)(E_\beta), \quad \beta = 1, \ldots, N.$$ 

3. the first nontrivial coefficient of its expansion in $t_0$ at $A_0$ is normalised to one.

Existence and uniqueness of the function $\psi(n_1, n_2, n_3)$ is due to application of the Riemann–Roch theorem with general position assumption and due to normalisation. Moreover, the function $\psi(n_1, n_2, n_3)$ is $\mathbb{K}$-rational, which follows from $\mathbb{K}$-rationality conditions for sets of points in their definition.

The next step of the construction is to obtain linear equations for wave functions. The full form of such equation is in case when the pole of $\psi(n_1, n_2, n_3)$ at $A_0$ is of the order exactly $(n_1 + n_2 + n_3)$ and respective zeros at $A_i$ are of the order $n_i$, for $i = 1, 2, 3$. We will call this case generic. Having fixed $\mathbb{K}$-rational local parameters $t_i$ at $A_i$, $i = 1, 2, 3$, denote by $\zeta_k(i)(n_1, n_2, n_3)$, $i = 1, 2, 3$, the $\mathbb{K}$-rational coefficients of expansion of $\psi(n_1, n_2, n_3)$ at $A_i$, respectively, i.e.,

$$\psi(n_1, n_2, n_3) = t_i^{n_i} \sum_{k=0}^{\infty} \zeta_k(i)(n_1, n_2, n_3)t_i^k, \quad i = 1, 2, 3.$$ 

Denote by $T_i$ the operator of translation in the variable $n_i$, $i = 1, 2, 3$, for example, $T_2\psi(n_1, n_2, n_3) = \psi(n_1, n_2+1, n_3)$. The full linear equation is of the form

$$T_i \psi - T_j \psi + \frac{T_j s_0(i)}{\zeta_0(i)} \psi = 0, \quad i \neq j, \quad i, j = 1, 2, 3.$$
It follows from observation that $T_i\psi - T_j\psi \in L(D)$, hence it must be proportional to wave function $\psi$. Coefficients of proportionality can be obtained from comparison (the lowest degree terms) of expansions of left and right sides of (1) at the point $A_i$.

Remark. When the genericity assumption fails then the linear problem (1) degenerates to the form $T_i\psi = \psi$ or even to $0 = 0$.

Notice that equation (1) gives

$$\begin{align*}
T_j\zeta_0^{(i)} &= -\frac{T_i\zeta_0^{(j)}}{\zeta_0^{(i)}}, & i \neq j, \quad i, j = 1, 2, 3.
\end{align*}$$

Define

$$\rho_i = (-1)^{\sum_{j<i} n_j} \zeta_0^{(i)}, \quad i = 1, 2, 3,$$

then equation (2) implies existence of a $\mathbb{K}$-valued potential (the $\tau$-function) defined (up to a multiplicative constant) by formulas

$$\begin{align*}
\frac{T_i\tau}{\tau} &= \rho_i, \quad i = 1, 2, 3.
\end{align*}$$

Finally, by "cyclic" use of equations (1) and (2) one can get the condition

$$\begin{align*}
\frac{T_2\rho_1}{\rho_1} - \frac{T_3\rho_1}{\rho_1} + \frac{T_3\rho_2}{\rho_2} &= 0,
\end{align*}$$

which, written in terms of the $\tau$-function, gives the discrete KP equation (8) called also the Hirota equation

$$\begin{align*}
(T_1\tau) (T_2T_3\tau) - (T_2\tau) (T_3T_1\tau) + (T_3\tau) (T_1T_2\tau) &= 0.
\end{align*}$$

Remark. Equation (8) can be obtained also from expansion of equation (1) at $A_k$, where $k = 1, 2, 3, k \neq i, j$.

Absence of a term in the linear problem (1), due to the Remark above, reflects in absence of the corresponding term in equation (8). This implies that in the non-generic case, when we have not defined the $\tau$-function yet, we are forced to put it to zero.

2.2. Explicit formulas for projective line. Denote by $t$ standard parameter on projective line $\mathbb{P}(\mathbb{K})$ and choose $A_0 = \infty$. Then we have the explicit formulas: vacuum wave functions

$$\psi^0 = (t - A_1)^{n_1} (t - A_2)^{n_2} (t - A_3)^{n_3}$$

and vacuum $\tau$-functions

$$\tau^0 = (A_1 - A_2)^{n_1n_2} (A_1 - A_3)^{n_1n_3} (A_2 - A_3)^{n_2n_3}.$$
Non vacuum $\tau$-function can be constructed by (see \[6\])

\[\tau = \tau^0 \det \phi_A^0(D, E),\]

where we denote by $\phi_A^0(D, E)$ the $N \times N$ matrix whose element in row $\beta$ and column $\alpha$ is

\[\phi_A^0(D, E)_{\alpha \beta} = \phi_\alpha^0(D_\beta) - \phi_\alpha^0(E_\beta), \quad \alpha, \beta = 1, \ldots, N.\]

and auxiliary vacuum wave functions $\phi_\alpha^0, \alpha = 1, \ldots, N$ have the form

\[\phi_\alpha^0 = \frac{1}{t - C_\alpha} \cdot \frac{(t - A_1)^{n_1}(t - A_2)^{n_2}(t - A_3)^{n_3}}{(C_\alpha - A_1)^{n_1}(C_\alpha - A_2)^{n_2}(C_\alpha - A_3)^{n_3}}.\]

**Remark.** Let $L = \mathbb{F}_q'$ be a field of rationality of all the points used in the construction. Then $\tau$-function is periodic in $n_1$, $n_2$ and $n_3$ with periods being divisors of $q^l - 1$, which is the order of the multiplicative group $\mathbb{L}_\ast$.

Algebro-geometric solutions can be obtained using the Jacobian of the curve $C$. Construction of KP cellular automata from hyperelliptic curves was presented in [4] and of KdV cellular automata in [2].

3. **Reduction to the discrete 1D Toda equation**

3.1. **Reduction to the discrete 1D Toda equation.** The discrete 1D Toda equation

\[(T_1 T_3^{-1} \tau) (T_1^{-1} T_3 \tau) - (T_1 \tau) (T_1^{-1} \tau) + \tau \tau = 0.\]

can be obtained from the dKP equation

\[(T_1 \tau) (T_2 T_3 \tau) - (T_2 \tau) (T_3 T_1 \tau) + (T_3 \tau) (T_1 T_2 \tau) = 0.\]

by imposing the constraint

\[T_1 T_2 T_3^{-1} \tau = \gamma \tau,\]

where $\gamma$ is a non-zero constant. The reduction \[8\] can be expressed at the level of the wave function $\psi$ as follows.

**Lemma.** Assume that on the algebraic curve $C$ there exists a rational function $h$ with the following properties

1. the divisor of the function is $(h) = -A_0 + A_1 + A_2 - A_3$,
2. it satisfies $N$ constraints: $h(D_\beta) = h(E_\beta), \beta = 1, \ldots, N$,
3. the first nontrivial coefficient of its expansion in the parameter $t_0$ at $A_0$ is normalised to one.

Then the wave function $\psi$ satisfies the following condition

\[T_1 T_2 T_3^{-1} \psi = h \psi.\]
Proof. Orders of poles and zeros of functions from both sides of equation (3.1) are the same, so these functions are proportional. Normalisation conditions (item 3 in Lemma and item 3 in Definition) imply they are equal. □

Remark. Equation (7) is identical with a discrete version of sine-Gordon equation considered in [10]. It is different from the standard form obtained by Hirota. As a consequence, cellular automata proposed here are different from those investigated in [5].

3.2. A \( \tau \)-function for the discrete 1D Toda equation. Using reduction condition for the wave function \( \psi \) we can derive \( \tau \)-function satisfying the discrete 1D Toda equation.

Proposition. Let \( h \) be the function as in Lemma 3.1. Let \( \delta_1, \delta_2 \) and \( \delta_3 \) denote the respective first coefficients of local expansions of \( h \) in parameters \( t_1, t_2 \) and \( t_3 \) at \( A_1, A_2 \) and \( A_3 \),

\[
h = t_1(\delta_1 + \ldots), \quad h = t_2(\delta_2 + \ldots), \quad h = \frac{1}{t_3}(\delta_3 + \ldots).
\]

Then the function

\[
\tilde{\tau} = \tau \delta_1^{-(n_1-1)/2}(-\delta_2)^{-n_2(n_2-1)/2} \delta_3^{n_3(n_3-1)/2}
\]

satisfies the discrete 1D Toda equation (7).

Proof. Expansions of equation (3.1) at points \( A_1, A_2 \) and \( A_3 \) give

\[
T_1T_2T_3^{-1} \rho_1 = \delta_1 \rho_1, \quad T_1T_2T_3^{-1} \rho_2 = -\delta_2 \rho_2, \quad T_1T_2T_3^{-1} \rho_3 = \delta_3 \rho_3.
\]

The functions

\[
\tilde{\rho}_1 = \delta_1^{-n_1} \rho_1, \quad \tilde{\rho}_2 = (-\delta_2)^{-n_2} \rho_2, \quad \tilde{\rho}_3 = \delta_3^{n_3} \rho_3,
\]

satisfy

\[
T_1T_2T_3^{-1} \tilde{\rho}_i = \tilde{\rho}_i, \quad i = 1, 2, 3.
\]

These functions define new potential \( \tilde{\tau} \) related to \( \tau \) by (9). The new \( \tilde{\tau} \)-function satisfies dKP equation and due to condition (10) it fulfils condition (8). As a consequence \( \tilde{\tau} \) satisfies equation (7). □

3.3. Explicit formulas for genus \( g = 0 \). In the case of projective line the function \( h \) reads

\[
h(t) = \frac{(t - A_1)(t - A_2)}{(t - A_3)}.
\]

Then,

\[
\delta_1 = \frac{(A_1 - A_2)}{(A_1 - A_3)}, \quad \delta_2 = \frac{(A_2 - A_1)}{(A_2 - A_3)}, \quad \delta_3 = (A_3 - A_1)(A_3 - A_2).
\]
The formula (9) and the determinant formula (6) allow to find pure $N$-soliton solutions of the discrete 1D Toda equation over finite fields.

4. Examples

Example 1. The vacuum solution of the discrete 1D Toda in $\mathbb{F}_5$

First, in Figure 1 we present 0-soliton $\tilde{\tau}$-function taking values in the finite field $\mathbb{F}_5$. We fix parameters of the solution from $\mathbb{F}_5$ as follows: $A_1 = 0$, $A_2 = 2$, $A_3 = 3$. For this example the Remark at the end of section 2.2 and formula (9) imply relations $T_1^4 \tilde{\tau} = (\delta_1)^{-2} \tilde{\tau}$ and $T_3^4 \tilde{\tau} = (\delta_3)^2 \tilde{\tau}$. Because $\delta_1 = 4$ and $\delta_2 = 3$ the period in $n_1$ variable is equal to 4, and in $n_3$ variable is 8. It is also evident from the picture that the reduction constraint $T_2 \tilde{\tau} = \gamma T_1^{-1} T_3 \tilde{\tau}$ where $\gamma = 2$ is satisfied. Notice also there are no values 0 in a vacuum solution.

Elements of $\mathbb{F}_5$ are represented by:

-0, -1, -2, -3, -4.

![Vacuum solution](image)

**Figure 1.** The vacuum solution $\tilde{\tau}(n_1, n_2, n_3)$ of the discrete 1D Toda equation for $n_2 = 0$ and 1. Axes: $n_1$ directed to the right (from 0 to 14), $n_3$ directed upward (from 0 to 14).

Example 2. The breather solution of the discrete 1D Toda in $\mathbb{F}_5$

The next example, presented in the Figure 2, is an analog of a breather solution taking values in the finite field $\mathbb{F}_5$. Parameters of the solution are chosen in a ”symmetric” way from $\mathbb{F}_5^2 = \mathbb{F}_5[x]/(x^2 + x + 1)$ i.e. a quadratic algebraic extension of the field $\mathbb{F}_5$ by the polynomial $(x^2 + x + 1)$. An element $ax + b \in \mathbb{F}_5^2$ is denoted by $(ab)$. The adequate
Frobenius automorphism is denoted by $\sigma_F$. Parameters are fixed as follows:

$A_1 = (00), A_2 = (02), A_3 = (03); C_1 = (10), C_2 = \sigma_F(C_1) = (44); D_1 = (12), E_1 = \sigma_F(D_1) = (41), D_2 = (24), E_2 = \sigma_F(D_2) = (32);$

The $\tilde{\tau}$-function is normalised by $\tilde{\tau}(n_1 = 1, n_2 = 0, n_3 = 1) = 1$. Respective periods in variables $n_1$ and $n_3$ are equal to 12 and 24. The reduction constraint $T_2\tilde{\tau} = \gamma T_1^{-1} T_3\tilde{\tau}$ is satisfied (for $\gamma = 2$) as visible in the figure.

Elements of $\mathbb{F}_5$ are represented like in the previous example by:

- Red: $(00)$
- Yellow: $(01)$
- Green: $(02)$
- Cyan: $(03)$
- Purple: $(04)$

Figure 2. The breather solution $\tilde{\tau}(n_1, n_2, n_3)$ of the discrete 1D Toda equation for $n_2 = 0$ and 1. Axes: $n_1$ directed to the right (from 0 to 29), $n_3$ directed upward (from 0 to 29).

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References

[1] E. D. Belokolos, A. I. Bobenko, V. Z. Enol’skii, A. R. Its, and V. B. Matveev, *Algebro-geometric approach to nonlinear integrable equations*, Springer-Verlag, Berlin, 1994.
[2] M. Bialycki, *Integrable KP and KdV cellular automata out of a hyperelliptic curve*, Glasgow Math. J. **47A** (2005), 33–44.

[3] M. Bialycki and A. Doliwa, *The discrete KP and KdV equations over finite fields*, Theor. Math. Phys. **137**(1) (2003), 1412–1418.

[4] M. Bialycki and A. Doliwa, *Algebro-geometric solution of the dKP equation over a finite field out of a hyperelliptic curve*, Commun. Math. Phys. **253** (2005), 157–170.

[5] A. Bobenko, M. Bordemann, Ch. Gunn, and U. Pinkall, *On two integrable cellular automata*, Comm. Math. Phys. **158** (1993), 127–134.

[6] A. Doliwa, M. Bialycki, and P. Klimczewski, *The Hirota equation over finite fields: algebro-geometric approach and multisoliton solutions*, J. Phys. A: Math. Gen. **36** (2003), 4827–4839.

[7] D. M. Goldschmidt, *Algebraic functions and projective curves*, Springer, New York, 2003.

[8] R. Hirota, *Discrete analogue of a generalized Toda equation*, J. Phys. Soc. Jpn. **50** (1981), 3785–3791.

[9] I. M. Krichever, *Algebraic curves and nonlinear difference equations*, Uspekhi Mat. Nauk **33**:4 (1978), 215–216.

[10] I. M. Krichever, P. Wiegmann, and A. Zabrodin, *Elliptic solutions to difference non-linear equations and related many body problems*, Commun. Math. Phys. **193** (1998), 373–396.

[11] R. Lidl and H. Niederreiter, *Introduction to finite fields and their applications*, Univ. Press, Cambridge, 1994.

[12] H. Stichtenoth, *Algebraic function fields and codes*, Springer-Verlag, Berlin, 1993.

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