A Fast Algorithm for Generating Long Self-Affine Profiles

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We introduce a fast algorithm for generating long self-affine profiles. The algorithm, which is based on the fast wavelet transform, is faster than the conventional Fourier filtering algorithm. In addition to increased performance for large systems, the algorithm, named the wavelet filtering algorithm, a priori gives rise to profiles for which the long-range correlation extends throughout the entire system independently of the length scale.

I. INTRODUCTION

With the advent of the computer as a serious research tool, there has been a revolution in the quantitative description of processes and structures that before were deemed too complex. Two of the key concepts used for this description are the fractal and its close relative, the self-affine structure \cite{1}. In the early eighties, much effort was spent identifying and describing various physical systems having fractal or self-affine structure. As time went by, focus slowly shifted from pure description to asking why such structures would appear. This lead to the development of the science of complex growth phenomena. Now, many aspects of these are well understood. However, there are still hosts of interesting but unanswered questions lingering on — see e.g. Refs. \cite{2,3} for recent reviews. More recently, focus has again begun to shift somewhat, and one sees work dealing with the physical consequences of the presence of fractal or self-affine structures. A concrete example of these three levels of development may be found in the study of fracture surfaces. In the early eighties, Mandelbrot et al. \cite{5} characterized fracture surfaces as self-affine, in the early nineties attempts were made to understand why fracture surfaces are self-affine \cite{6}. Recently, phenomena such as two-phase flow in fracture joints have been studied \cite{7}.

In order to study the physical consequences of the presence of self-affine surfaces, algorithms generating these must be found. There are already several in existence, see e.g. Feder \cite{1}. However, subtle phenomena require the generation of huge surfaces. Two aspects of the algorithms then become important: (i) how self-affine are the surfaces that are generated, and (ii) how fast is the algorithm.

The most popular algorithm used today is the Fourier filtering algorithm. This algorithm, which has a fast implementation thanks to the fast Fourier transform, consists of generating in the Fourier domain, uncorrelated Gaussian random numbers which are filtered by a decaying power-law filter of exponent $-2H - 1$ where $H$ is the Hurst exponent (to be defined in Section II). By taking advantage of the inverse fast Fourier transform, self-affine surfaces in real space, with the desirable correlations, are generated.

It has previously been shown \cite{8} that the Fourier filtering algorithm has the disadvantage that the self-affine correlations in the limit of large systems only exists over a fraction of the total system size. This is due to aliasing effects. For large enough systems this fraction might be well below 1%. One might overcome the above problem by e.g. temporally generating a much larger surface than actually needed, and only use a small fraction of the total size. This is, however, not a very appealing approach, as the computer time and memory needed easily become too large. Another way of getting around this problem is due to Makse et al. \cite{10}. Here the (Fourier-space) filter function is modified by the introduction of a large momentum cut-off through the use of a modified Bessel function in the Fourier transform of the power spectrum. They show that this large momentum cut-off, while irrelevant for the large scale behavior in real space, is essential in order to suppress the aliasing effect and thereby obtaining sur-
faces with the desired scaling properties over the entire system size.

In this paper we report on an alternative filtering algorithm based on wavelets which a priori, and without modifications, gives rise to self-affine correlations which extends (up to finite size effects) over the entire system independently of its size. This algorithm is also computationally cheaper than the traditional (or modified) Fourier filtering algorithm.

This paper is organized as follows: In Section II we briefly review the defining properties of self-affine surfaces. Here we also include some results which will prove useful to us later. Section III is devoted to the outline of the new algorithm. In Section IV we present numerical studies of the this algorithm. We conclude in Section V.

II. SELF-AFFINE SURFACES

We limit our discussion in this paper to (1 + 1)-dimensional surfaces, which we will call profiles. A (statistically) self-affine profile, \( h(x) \), is by definition a structure which remains (statistically) invariant under the following scaling relation

\begin{align}
  x & \rightarrow \lambda x, \\
  h & \rightarrow \lambda^H h.
\end{align}

Here \( \lambda \) is a real number and \( H \), known as the roughness or Hurst exponent, characterizes this invariance. This exponent is usually in the range from zero to one. When \( H = 1/2 \), the profile is not correlated. An example of such a profile is the Brownian motion in one dimension. In this case, we interpret time as \( x \) and \( h(x) \) as the position on the Brownian particle at time \( x \). When \( H > 1/2 \) the profile is persistent, while when \( H < 1/2 \) it is anti persistent.

We show in Fig. 1 an example of a self-affine profile generated by the algorithm to be presented in this paper.

From the scaling relation (2.1), one can often relatively easily derive scaling relations for related quantities. In this paper we will later explicitly need the scaling relation for the second order structure function

\[ S(\Delta x) = \left\langle |h(\Delta x + x) - h(x)|^2 \right\rangle_x, \]

where \( \langle \cdot \rangle_x \) represents the average over the position variable \( x \), and the power spectrum \( P(q) \), defined as the Fourier transform of the height-height correlation function. They scale as

\[ S(\Delta x) \sim (\Delta x)^{2H}, \]

and

\[ P(q) \sim q^{-2H-1}. \]

We will make use of these two scaling relations in Section V.

III. THE WAVELET FILTERING ALGORITHM

Recently, the wavelet transform (14,15] has been used to analyze self-affine profiles (14,15]. The idea behind this analysis is as follows: We denote the wavelet transform of a function \( h(x) \) by \( W[h](a, b) \), where \( a \) and \( b \) are the scaling and location variables respectively. They form together the wavelet domain. In Ref. 15 the authors introduced what they called the average wavelet coefficient function, defined as \( W[h](a) = \langle |W[h](a, b)| \rangle_b \), where \( \langle \cdot \rangle_b \) denotes the average over all the location parameters \( b \) corresponding to one and the same scale \( a \). For a self-affine function, \( h(x) \), this quantity should scale as

\[ W[h](a) \sim a^{H+1/2}. \]

In much the same way as the Fourier filtering algorithm is used for generating self-affine profiles via the fast Fourier transform, a wavelet based filtering technique can be based on Eq. (3.1) in combination with the fast wavelet transform. The output of the fast wavelet transform (14,15] is a vector organized as a collection of various levels or hierarchies all of different lengths where each level, \( \ell \), is associated with a corresponding scale \( a_\ell \). The two first components of this vector, also known as level \( \ell = 0 \), are associated with the scaling function. All the other components are “true” wavelet coefficients, such that at level \( \ell \), corresponding to scale \( a_\ell = 2^{-\ell} \), there are \( N_\ell = 2^\ell \) coefficients. These coefficients (using our convention) are arranged such that the coefficients of the highest level are found at the end of the vector, and the levels decrease monotonically towards the top of the vector, corresponding to the level \( \ell = 0 \).

Hence, the wavelet based algorithm which we will be referring to as the wavelet filtering algorithm (WFA), consists of the following three steps:

1. Strictly speaking, one should in order to fully define the self-affine structure, also introduce the topothesy which is defined as the length-scale \( l \), over which the RMS-height, \( \sigma \), is \( \sigma(l) = l \). However, we will not explicitly need this latter quantity here, and will therefore simply neglect any further reference to it.
• Generate in the wavelet-domain normalized uncorrelated Gaussian numbers \( \{ \eta_i \} \), with \( i = 1, 2, \ldots, N \), where \( N \) is the number of discrete points \( x_i \) that together with \( h_i = h(x_i) \) constitute the self-affine profile. \( N \) is assumed, due to the use of the fast wavelet transform, to be a power of 2.

• Filter these random numbers according to

\[
w_i = \left( a_{\ell(i)} \right)^{H+1/2} \frac{\eta_i}{\langle |\eta| \rangle_{\ell(i)}}, \quad i = 1, 2, \ldots, N,
\]


\( w_i \) to obtain the wavelet coefficients \( \{ w_i \} \). Here \( a_{\ell(i)} = 2^{-\ell(i)} \) represents the scale, at level \( \ell(i) \), where \( \ell(i) \) is defined as the level corresponding to the location index \( i \) of the vector \( u_i \) (or \( \eta_i \)). Furthermore, \( \langle |\eta| \rangle_{\ell(i)} \) represents the average of the absolute value of those \( \eta_i \) that together form level \( \ell(i) \).

• Perform the inverse fast wavelet transform on \( \{ w_i \} \), with the (compactly supported) wavelet of your choice, to obtain the (real-space) self-affine profile of predefined Hurst exponent \( H \).

With the wavelet filtering algorithm, good quality self-affine surfaces with predefined Hurst exponent can be generated. In Fig. 1 we show an example of a self-affine surface with predefined Hurst exponent \( H = 0.6 \) and length \( N = 4096 \) generated by the algorithm just outlined. It is worth noting that the above three steps can be modified in order to deal with surfaces in higher dimensions [16]. In this case the speed of the surface generating algorithm becomes very important.

One of the prominent features of the wavelet transform is that the basis functions, the wavelets, are localized in both space and frequency. This has as a consequence, among others, that there is no aliasing, or at least heavily suppressed as compared to the Fourier transform. This implies that the wavelet filtering algorithm should automatically result in surfaces which have the desired correlations over the entire length of the profile, and not just a small fraction of it. Hence, independent of the system size, the WFA is capable of generating profiles with well-defined long-range correlations. We will demonstrate the validity of this claim in Section IV.

Before turning to the numerical studies of the WFA, we add some remarks regarding the computational efficiency of this algorithm. The most time-consuming part of the WFA is the inverse wavelet transform. To a good approximation, at least for larger system sizes, this time determines the overall computational time of the entire algorithm. The fast wavelet transform needs \( \mathcal{O}(cN) \) operations, where \( c \) is a positive real number which value depends on the wavelet used [13-17]. Thus, the number of operations needed for generating a surface by WFA is linear in the number of points belonging to the profile. In comparison, the Fourier filtering algorithm, which speed is mainly controlled by the fast Fourier transform, needs \( \mathcal{O}(N \log_2 N) \) to generate a profile. For large system sizes the difference in execution time between WFA and the Fourier filtering algorithm becomes significant.

IV. NUMERICAL RESULTS

In order to test numerically the predictions of the previous section, we have chosen to study the second order structure function, \( S(\Delta x) \), and the power spectrum \( P(q) \) of self-affine profiles generated with the wavelet filtering algorithm. The appropriate scaling relations for these two quantities are given by Eqs. (2.3) and (2.4). They will provide us with independent information enabling us to accurately quantify over which length scales the self-affine correlations exist. The numerical experiments, for which the results will be presented shortly, were performed as follows: We generated, by WFA, an ensemble of long self-affine profiles all with the same Hurst exponent \( H \). For each profile the structure function and power spectrum were calculated, and these were averaged over the ensemble of profiles.

In Fig. 2 we give the numerical results for the second order structure function obtained as described above. The predefined Hurst exponents, used by the surface generator, were from bottom to top \( H = 0.8, 0.6, 0.4, \) and 0.2 as indicated in the figure. The length of each profile was \( N = 2^{25} = 33554432 \). The number of profiles used in obtaining the averages was \( N_h = 50 \), and the wavelet used was of the Daubechies-type (D12) [12-13]. The dashed lines are regression fits to the numerical data. They correspond (from bottom to top) to Hurst exponents of \( H = 0.80 \pm 0.01, 0.60 \pm 0.01, 0.41 \pm 0.01, \) and \( 0.22 \pm 0.02 \), all consistent, within the errorbars, with the predefined exponents given above. One easily observes from Fig. 2 that the correlations extend over all scales except for the largest lags \( \Delta x \). The reason that the last few large lags do not fit into this general picture is due to finite size effects. We have also undertaken the above analysis for different types of wavelets, taken from the Daubechies family, and for various system sizes, finding no results that are inconsistent with those presented in Fig. 2.

In order to make the comparison with the Fourier filtering algorithm even more apparent, we present in Fig. 3 the average power spectrum obtained using the same surfaces as were used in Fig. 2. The correlations again span most scales. The dashed regression fits lead to the following exponents (from bottom to top): \( H = 0.80 \pm 0.01, 0.61 \pm 0.01, 0.41 \pm 0.01, \) and \( 0.20 \pm 0.01 \), which again is in excellent agreement with the values of the Hurst exponent used for the generation of the underlying profiles.

Figs. 2 and 3 indicate that the self-affine correlations span all but the largest scales of the profiles. We stress that this is a generic property of the wavelet filtering algorithm, and no modification of the algorithm is needed in order to handle large system sizes in a satisfactory manner. This is a consequence of the celebrated property of

\[ \Delta x = 0 \]
the wavelets being localized both in space and frequency.

The calculations of this paper were performed on a SGI/Cray Origin 2000 supercomputer based on the R10000 chip from SGI. On this machine the average CPU time needed for generating a profile of the length used above \( N = 2^{25} \) was \( t_{\text{WFA}} = 45s \) and \( t_{\text{FFM}} = 125s \) for the WFA and the traditional (or modified) Fourier filtering algorithm respectively. Hence the speedup gained by using the wavelet filtering algorithm over the Fourier filter algorithm is close to a factor 3. For system sizes \( N \sim 10^3 \), we could not observe any significant different between the two algorithms.

V. CONCLUSIONS

To summarize, we have introduced a fast and simple algorithm for generating long (or short) self-affine profiles. This algorithm, named the wavelet filtering algorithm, is demonstrated to overcome the problem related to the aliasing effect which the traditional Fourier filtering algorithm is troubled with. Furthermore, the wavelet based filtering technique outperforms its Fourier-domain counterpart by large margins with respect to computational costs, at least for large system sizes.

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FIG. 1. A self-affine profile, \( h(x) \), of length \( N = 4096 \) and Hurst exponent \( H = 0.6 \) generated by the wavelet filtering algorithm. The wavelet used in order to generate the profile was the D12 Daubechies wavelet.

FIG. 2. The average second order structure function, \( S(\Delta x) \), obtained by averaging over \( N_h = 50 \) samples (for given \( H \)) of the self-affine profile \( h(x) \) generated by WFA. All profiles were of length \( N = 2^{25} = 33554432 \). The Hurst exponents used were (from bottom to top) \( H = 0.8, 0.6, 0.4, \) and \( 0.2 \), as indicated in the figure. The dashed lines are the best regression fits corresponding respectively to Hurst exponents (from bottom to top) \( H = 0.80 \pm 0.01, 0.60 \pm 0.01, 0.41 \pm 0.01, \) and \( 0.22 \pm 0.02 \).

FIG. 3. The average power spectrum \( P(q) \) obtained by averaging over \( N_h = 50 \) samples (for given \( H \)) of the self-affine profile \( h(x) \) generated by WFA. All profiles were of length \( N = 2^{25} = 33554432 \). The Hurst exponents used were (from bottom to top) \( H = 0.8, 0.6, 0.4, \) and \( 0.2 \), as indicated in the figure. The dashed lines are the best regression fits corresponding respectively to Hurst exponents (from bottom to top) \( H = 0.80 \pm 0.01, 0.61 \pm 0.01, 0.41 \pm 0.01, \) and \( 0.20 \pm 0.01 \).
Figure 1
I. Simonsen and A. Hansen
A new fast algorithm for generating long self-affine profiles
Figure 2
I. Simonsen and A. Hansen
*A new fast algorithm for generating long self-affine profiles*
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