Investigation into microstructure and element segregation of Inconel 713LC during SLM processing using modified cellular automaton model

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Abstract

A numerical and experimental investigation is performed into the feasibility of fabricating IN713LC components using selective laser melting (SLM) through an appropriate control of the solidification speed, microstructure formation and element segregation. A modified Cellular Automaton (CA) model is developed to explore the nucleation, grain growth and element segregation behavior of IN713LC during ultra-fast solidification. It is found that the undesired phase formation which occurs during SLM processing of IN713LC is caused by the micro segregation of Nb, Ti, Mo and C at the grain boundaries. It is further shown that the micro segregation intensity depends on the solidification speed, which is determined in turn by the laser energy density. In particular, a lower laser energy density increases the solidification speed and results in a more uniform solid phase, thereby reducing the risk of crack formation. The simulation results are verified by experimental investigation. The results confirm that a lower laser energy density reduces the crack density and crack length. Finally, a crack-free IN713LC SLM sample is successfully produced by reducing the energy density from 360 to 170 J/m.

Introduction

IN713LC, a nickel alloy with a composition of Cr-Al-Mo-Ti-Nb-Zr, is known for its good fatigue resistance, excellent mechanical properties, superior oxidation resistance and enhanced resistance to degradation under harsh operating condition1–3. Many studies have shown that Ni-based superalloys, such as IN7184,5, IN6256 and Hastelloy X7 can be successfully fabricated by selective laser melting (SLM). However, due to its crack-prone nature, the SLM processing of IN713LC is extremely challenging. Previous studies have reported the formation of three main phases in IN713LC components built by SLM, namely γ, γ’ and carbide8,9. Among these phases, γ phase is the main phase and has the same composition as the nominal IN713LC alloy, while γ’ and carbides are secondary phases and have compositions of Ni3Nb and MC, respectively10, where M is the metal element and C is carbon.

Besides the usual factors affecting the processability of SLM, such as the thermal stress, raw powder material quality and pre-printing procedure, IN713LC is also prone to crack
formation as the result of a significant segregation and depletion of the strengthening elements from the γ phase during solidification\textsuperscript{1,11-13}. The formation of a large mushy zone further prompts crack formation through the suppression of liquid metal feeding in solidification zone\textsuperscript{14}.

Due to the extremely rapid rate of solidification during SLM, experimental investigations provide only limited insights into the microstructure formation and element segregation behavior during cooling. Thus, in exploring the mechanisms which render the SLM processing of IN713LC so challenging, simulation models are widely preferred. Cellular Automaton (CA) is recognized as one of the best simulation methods for examining ultra-fast solidification processes such as SLM. Ao et al\textsuperscript{15} used CA to simulate the microstructure of AlSi10Mg during SLM based on an analysis of the thermodynamic-driven growth kinetics of the solid-liquid interface and melt pool composition. Trivedi\textsuperscript{16} examined the grain growth rate using the Kurz-Giovanola-Trivedi model. Their simulation demonstrated that the cooling rate during SLM process was around $10^5 \text{ to } 10^6$ K/s. It was additionally shown that during the solidification process, the equiaxed grain formation increased with an increasing pre-heating temperature or reducing scanning speed. Zinoviev et al\textsuperscript{17} developed a two-dimensional CA model to evaluate the grain structure of metallic alloys during SLM processing. The results showed that the CA model yielded accurate predictions of the main characteristics of the alloy microstructure following SLM. Lain et al\textsuperscript{18} examined the microstructure evolution of IN718 during additive manufacturing using a three-dimensional CA model and showed that the CA model can simulate the different microstructure forms of IN718 such as columnar grains, fine and course equiaxed grains using a complex 3D transition thermal model.

However, the literature presently lacks any simulation studies on the grain structure evolution and element segregation of IN713LC during SLM processing. Accordingly, the current study develops a modified CA model to predict the microstructure and microsegregation behavior of IN713LC during rapid solidification. The model takes account of boundary and bulk nucleation and grain growth and equations of solute diffusion have been also coupled with CA model. A particular effort has been made to develop a practical and accurate model for elements partition coefficient during SLM processing under conditions far from equilibrium conditions. The simulation results are used to determine the SLM conditions which enable the crack-free printing of IN713LC components. The validity of the derived conditions is demonstrated by means of experimental trials.

Results

The microstructures and element compositions of the IN713LC samples fabricated using energy densities of 360, 280 and 210 J/m (corresponding to cooling rates of $4.3 \times 10^4$, $5.5 \times 10^4$ and $7.3 \times 10^4$ K/s, respectively) were simulated using the CA model. The simulations were conducted in 2D boxes with widths and heights of 200 µm × 90 µm, 175 µm × 80 µm and 150 µm × 70 µm, corresponding to energy densities of 360, 280 and 210 J/m, respectively. (Note that the box dimensions were set in accordance with the melt pool size observed in experimental single laser track SLM trials performed under the same energy density conditions.)

The composition of the initial IN713LC powder is shown in Table 1. In implementing the CA model, the IN713LC alloy was considered to be a 6-element system consisting of Cr-Mo-Nb-Al-Ti-C (i.e., the matrix component Ni was deliberately excluded). Note
that the C component of the alloy was retained in the model even though its concentration is extremely low (0.05 wt.%, see Table 1) due its importance in prompting hot cracking. Furthermore, equal diffusion was assumed for all of the components in both the solid phase and the liquid phase, and no-flux boundary conditions were applied at all the boundaries. Table 2 summarizes the parameters applied in the micro-segregation model for the different elements of the IN713LC alloy. It is noted that $C_0$ denotes the initial element concentration in the powder, and $k$ is the equilibrium partition coefficient obtained in the low-speed continuous casting of IN713LC. The physical parameters employed in the CA model are listed in Table 3.

Table 1 Chemical composition of IN713 LC (mass %)

|   | Ni  | Cr  | Mo  | Nb  | Al  | Ti  | C  |
|---|-----|-----|-----|-----|-----|-----|----|
| Bal.| 12.1| 4.1 | 2.0 | 6.2 | 0.77| 0.05|    |

Table 2 Parameters used in micro-segregation model for different alloying elements of IN713LC.

| Parameter | Cr   | Mo   | Nb   | Al   | Ti   | C   |
|-----------|------|------|------|------|------|-----|
| $C_0 \times \% \quad [\text{mass}]$ | 12.1 | 4.10 | 2.0  | 6.2  | 0.77 | 0.05|
| $D \times 10^{13} \quad [\text{m}^2/\text{s}]$ | 8.98 | 10.76| 10.53| 11.10| 10.99| 43.0|
| $0.96$ | 0.82 | 0.46 | 0.83 | 0.55 | 0.12 |

Table 3 Physical parameters used in CA model.

| Parameter       | Value |
|-----------------|-------|
| CA Cell size[µm]| 0.2   |
| $n_{x,max}[1/\text{m}^3]$ | $1.5 \times 10^6$ |
| $n_{v,max}[1/\text{m}^3]$ | $1.5 \times 10^7$ |
| $T_L[K]$        | 1622  |
| $T_S[K]$        | 1523  |
| $\Delta T_{S,\max}[K]$ | 3.5  |
| $\Delta T_{S,\sigma}[K]$ | 0.2   |
| $\alpha[\text{m}^2/\text{s}]$ | $4.7 \times 10^6$ |
| $k_f[\text{J/(s m K)}]$ | 18.3  |
| $\Gamma[K/\text{m}]$ | $3.15 \times 10^7$ |
| $\Delta H[\text{kJ/kg}]$ | 236   |
| $c_p[\text{J/(g·K)}]$ | 0.668 |

Table 4 lists the results obtained of the CA model for the non-equilibrium partition coefficients of Cr, Mo, Nb, Al, Ti and C during solidification given different values of the laser energy density. It is seen that the different elements have different solubilities in the IN713LC matrix. In particular, the C, Nb, Ti, Al and Mo elements have relatively low solubilities in the solid matrix, and hence tend to segregate in liquid. By contrast, Cr has a higher solubility, and thus tends to be trapped in the solid. It is noted that C has the highest diffusion coefficient of all the components in IN713LC and therefore moves through the solid-liquid interface easier than the other elements. The rapid diffusion of C atoms can be attributed to their small size, which allows them to sit interstitially within the Ni-super alloy lattice and hence move faster by jumping from one interstice to another.
Comparing the equilibrium partition coefficients in Table 2 with the non-equilibrium element partition coefficients in Table 4, it is seen that their values tend toward unity as the cooling rate increases, i.e., the energy density reduces. In other words, a more uniform element concentration is obtained under lower energy densities. This finding is reasonable since, as the energy density decreases, the solidification speed increases, and hence the elements have less time to diffuse from the solid phase to the liquid phase. In other words, the high-speed solid-liquid interface traps the elements in the solid phase before they can pass through the solid-liquid interface and reach the liquid phase. It is hence inferred that different segregated components will be formed under higher solidification speeds.

Table 4. Average IN713LC element partition coefficients under different solidification rates.

| Energy Density [J/m] | Cr   | Mo   | Nb   | Al   | Ti   | C    |
|----------------------|------|------|------|------|------|------|
| 360                  | 0.97 | 0.87 | 0.60 | 0.84 | 0.71 | 0.15 |
| 280                  | 0.98 | 0.89 | 0.63 | 0.85 | 0.73 | 0.17 |
| 210                  | 0.98 | 0.90 | 0.71 | 0.86 | 0.76 | 0.21 |

Figures 1~3 show the simulation results for the grain structure and Nb, Ti and C concentrations at different solidification percentages (i.e., 10%, 50% and 100%) and energy densities of 360, 280 and 210 J/m, respectively. (Note that Nb, Ti and C are deliberately selected here since they have the lowest partition coefficients among all the elements in IN713LC, as shown in Table 4.) In general, the results reveal that a lower energy density leads to a more dominant columnar grain growth structure along the solidification direction. In addition, a higher energy density results in a greater number of equiaxed grains above the columnar grains (see Fig. 1, for example, for the maximum energy density of 360 J/m). The greater volume fraction of equiaxed grains can be attributed to the lower temperature gradient produced at higher energy densities (i.e., 4.5×10⁵ K/m at an energy density of 360 J/m, compared to 5.8×10⁵ and 7.8×10⁵ K/m at energy densities of 280 J/m and 210 J/m, respectively). More specifically, a lower temperature gradient results in the formation of equiaxed grains in front of the columnar grains, whereas a higher temperature gradient promotes directional growth.

The element concentration distribution results shown in Figs. 1~3 show that the surface and grain boundaries are enriched in Nb, Ti and C. A high element concentration at the surface is reasonable since the surface is the last remaining liquid region during the solidification process, and is enriched with segregated elements from the lower stream of the melt pool. Once this region solidifies, the elements have no other liquid phase to be rejected to and are hence confined to the solid phase at the surface. The element distribution results in Fig. 1 show an obvious segregation of the Nb, Ti and C elements at the grain boundaries, even under a 50% solidification state. However, no such micro-segregation effect is observed under lower energy densities of 280 J/m and 210 J/m, respectively. Hence, it can be inferred that SLM processes performed under higher energy densities (i.e., slower solidification speeds) will produce more intensive micro segregation of the elements. In other words, the likelihood of carbide or secondary phase formation at the grain boundaries is increased specially with the present of Nb and Ti. Nb is founded as a strong carbide former. In addition, based on the previous research, Ti is also known to control the formation of carbide during the solidification. Then, segregation of both Nb and Ti during the solidification is important and even small difference may result in different conditions.
Figure 1. Simulation results for solidification behavior at energy density of 360 J/m. (a) grain structure after 10% solidification, (b) grain structure after 50% solidification, (c) grain structure after full solidification. (d) Nb concentration after 10% solidification, (e) Nb concentration after 50% solidification, (f) Nb concentration after full solidification. (g) Ti concentration after 10% solidification, (h) Ti concentration after 50% solidification, (i) Ti concentration after full solidification. (j) C concentration after 10% solidification, (k) C concentration after 50% solidification, (l) C concentration after full solidification.

Figure 2. Simulation results for solidification behavior at 280 J/m. (a) grain structure after 10% solidification, (b) grain structure after 50% solidification, (c) grain structure after full solidification. (d) Nb concentration after 10% solidification, (e) Nb concentration after 50% solidification, (f) Nb concentration after full solidification. (g) Ti concentration after 10% solidification, (h) Ti concentration after 50% solidification, (i) Ti concentration after full solidification. (j) C concentration after 10% solidification, (k) C concentration after 50% solidification, (l) C concentration after full solidification.
A detailed inspection of the simulation results presented in Figs. 1~3 shows that not all the grain boundaries exhibit the same amount of segregation. For example, when large and elongated grains encounter other grains during the growth process, a grain boundary with a high element concentration is formed. In general, grains with a larger size inevitably accrue a greater amount of segregation in front of their solid-liquid interface. When the segregated liquid is confined by other grains in the upstream region of the melt pool (particularly in the dense area of equiaxed grains), the segregated elements are confined at the interface between them, and hence a region of high element concentration is formed. Notably, however, the equiaxed grains may sometimes decrease the formation of segregated elements if, for example, they are formed directly in front of a columnar grain when the segregated level ahead of the solid-liquid interface of columnar grain still is not too high. The regions indicated by lines AA’ and BB’ in Figs. 1(c, f, j and l) are typical cases in point. For example, AA’ shows a uniform and dense equiaxed grain located immediately above the melt pool boundary. The presence of this grain physically suppresses element segregation at the boundary of the advancing grain. By contrast, line BB’ shows the presence of an equiaxed grain in front of a large grain with a high level of segregation at the solid front. In this condition, a high segregated area is formed within the system. Base on the role of equiaxed grains on segregation, many researchers have attempted to suppress element segregation at the grain boundary in additive manufacturing processes by lowering the temperature gradient in order to increase the number of equiaxed grains. However, pre-heating requires the use of a middle frequency inductive power to provide heat to the substrate and a thermometer to perform temperature monitoring. Consequently, the present study deliberately adopts a simpler strategy for controlling the grain boundary segregation effect in the SLM processing of IN713LC by controlling the solidification speed through an appropriate setting of the laser power and scanning speed, respectively.

Figure 4 presents optical microscope (OM) images of the ground and polished IN713LC samples printed using energy densities of 360, 280 and 210 J/m, respectively. It is seen that all of the samples contain cracks of various lengths. For each sample, the cracks are randomly dispersed and are orientated mainly parallel to the solidification
direction. The OM images also show the presence of sub-micrometric spherical pores formed as the result of gas entrapment\(^\text{24}\). As predicted by the simulation results, the crack density and crack length decrease with a reducing energy density. In particular, the crack density and average crack length reduce from 0.103% and 79 \(\mu\)m for an energy density of 360 J/m to 0.029% and 62 \(\mu\)m for an energy density of 210 J/m, as shown in Table 5. According to the simulation results, the lower degree of cracking at a lower energy density can be attributed to the faster cooling rate, which suppresses element segregation at the grain boundaries and generates a more uniform element distribution. As a result, the formation of undesired phase, which leads to crack initiation, is reduced.

Table 5. Crack density and average crack length for different energy densities.

| Energy density [J/m] | Crack density [%] | Average crack length [\(\mu\)m] |
|----------------------|-------------------|-------------------------------|
| 360                  | 0.103             | 79                           |
| 280                  | 0.082             | 69                           |
| 210                  | 0.029             | 62                           |

Figure 4 presents a high-magnification SEM image and corresponding EBSD results for the crack region of the IN713LC sample built using an energy density of 210 J/m. A dominant cracking pattern is clearly observed along the grain boundary region. The cracks extend along the grain boundaries, or are located at the point of intersection of neighboring grains. Overall, the straight shape of the cracks, and their location at the grain boundaries, indicates that they are formed as the result of segregation-induced grain boundary liquation\(^0\). Notably, the experimental results are consistent with the simulation predictions that the majority of the cracks are located at the grain boundaries.
To verify the role of element segregation in prompting crack formation in the present samples, a detailed EDS analysis was performed at selected points inside and outside of the cracks formed in the sample built with a high energy density of 360 J/m. The analysis points and corresponding EDS results are presented in Fig. 6. As predicted by the simulation results, both crack regions exhibit a high segregation rate of Nb, Ti, Al and C. In particular, the Ti, Nb, Al and C contents increase from 0.77, 2.0, 6.2 and 0.05 wt.% in the original IN713LC powder to 24.42, 11.54, 16.08 and 21.54 wt.%, respectively, within the crack. The high concentration of metallic elements implies the formation of MC carbides within the crack, and is hence consistent with the simulation results for the non-equilibrium element partition coefficients listed in Table 4.

In general, the simulation and experimental results presented above indicate that a lower energy density (i.e., a faster solidification rate) is beneficial in suppressing crack formation in SLM IN713LC components. Accordingly, a further sample was printed using a laser power of 150 W and a scanning speed of 850 mm/s; corresponding to a reduced energy density of around 170 J/m. To eliminate the sub-micrometric spherical pores formed at low energy densities, the scan spacing was reduced to compensate for the shrinkage of the melt pool. Figure 7 presents an OM image of the built sample. It is seen that the surface contains neither cracks nor visible pores. In other words, the feasibility of SLM for the processing of IN713LC is confirmed given a suitable low value of the laser energy density. It is noted, however, that reducing the energy density below a certain critical value may potentially induce new defects such as ‘void formation’ due to melt pool shrinkage under low energy conditions. Accordingly,
determining the optimal value of the energy density is an essential requirement for future studies.

Figure 7. IN713LC sample fabricated using low energy density of 170 J/m.

Discussion

A modified CA model has been developed to predict the microstructure and micro segregation behavior of IN713LC nickel alloy during the solidification stage of the SLM process. The model takes explicit account of the grain nucleation and growth mechanisms and solute diffusion of the alloying elements, respectively. The simulation results show that the micro segregation phenomenon is a diffusion-controlled process, in which the elements with a lower partition coefficient and higher diffusivity experience a higher rate of segregation. It is additionally shown that element segregation is enhanced under a slower solidification rate since the elements spend a longer time in their respective precipitation windows and thus have sufficient time to move from the solid phase to the liquid phase. The simulation results also indicate that, among all the alloying elements of IN713LC, the Nb, Ti, C elements exhibit the highest rate of micro segregation during solidification. Finally, as the laser energy density reduces, the solidification speed increases and the element partitioning coefficients approach unity. Consequently, a more uniform solid phase is formed with only minimal segregation at the grain boundaries. The simulation results have been verified by means of experimental trials. It has been shown that the crack length and crack density decrease with a reducing laser energy density. The experimental results are thus consistent with the simulation predictions for the effects of the laser energy density on the element segregation effect at the grain boundary during solidification. Notably, the experimental results have demonstrated for the first time the feasibility of fabricating crack-free IN713LC components using the SLM process given a suitable low value of the energy density (e.g., 170 J/m in the present case).
Methods

This section describes the modified CA model developed in the present study to examine the microstructural evolution of IN713LC during SLM processing and the corresponding element segregation behavior. The section commences by introducing the models used to describe the nucleation and growth of the IN713LC microstructure as a function of the undercooling temperature. The section then describes the approach taken to model the element segregation behavior based on the principles of solute diffusion. Finally, the section outlines the CA simulation procedure used to predict the grain structure and micro-segregation effect in IN713LC during the rapid solidification stage of the SLM process.

Nucleation and grain growth models

The number of nucleation sites formed during liquid phase solidification is conventionally assumed to have a Gaussian distribution as a function of the undercooling temperature and can be modelled using the following equation:

\[
\frac{dn}{d(\Delta T)} = \frac{n_{\text{max}}}{\sqrt{2\pi}} \exp \left[ -\frac{(\Delta T-\Delta T_{\text{max}})^2}{2\Delta T_{\sigma}^2} \right]
\]

(1)

where \(\Delta T\) is the undercooling temperature, \(\Delta T_{\text{max}}\) is the maximum nucleation undercooling temperature, \(\Delta T_{\sigma}\) is the standard deviation of the undercooling temperature, and \(n_{\text{max}}\) is the maximum nucleation density. The total density of the nucleation sites can thus be modeled as

\[
n(\Delta T) = \int_{0}^{\Delta T} \frac{dn}{d(\Delta T)} d(\Delta T).
\]

(2)

Since the IN713LC alloy system is dominated by Ni (~75wt.%, see Table 1), the undercooling temperature, \(\Delta T\), can be assumed to consist of only thermodynamic undercooling \(\Delta T_T\). That is,

\[
\Delta T = \Delta T_T.
\]

(3)

The thermodynamic undercooling, \(\Delta T_T\), can be expressed as:

\[
\Delta T_T = \frac{\Delta H}{c_p} \cdot \text{Iv}(P_t),
\]

(4)

where \(c_p\), \(\Delta H\), \(P_t\) and \(\text{Iv}\) are the specific heat, latent heat, Peclet number and Ivantsov function, respectively. The Ivantsov function has the form:

\[
\text{Iv}(P_t) = P_t \exp(P_t) \int_{P_t}^{\infty} \frac{e^{-\eta}}{\eta} d\eta.
\]

(5)

In addition, the Peclet number is defined as:

\[
P_t = \frac{\nu_{\text{growth}} R}{2\alpha},
\]

(6)

where \(\nu_{\text{growth}}\) is the liquid / solid growth speed, \(\alpha\) is the thermal diffusivity, and \(R\) is the dendrite tip radius. According to Trivedi26 and Kurz and Fisher27, the growth speed and dendrite tip radius are related via the following stability criterion:

\[
R = \frac{c_p \Gamma}{0.025 P_t \Delta H},
\]

(7)

where \(\Gamma\) is the Gibbs-Thomson coefficient. From Eqs. (6) and (7), the growth rate can be obtained as

\[
\nu_{\text{growth}} = \frac{2 P_t \Delta H}{d \nu}\]

(8)
Element segregation

In modeling the element segregation behavior during solidification, the governing equation for solute diffusion is given as:

\[
\frac{\partial C_i}{\partial t} = \nabla D_i(T) \nabla C_i + \left( C_{L,i}^* - C_{S,i}^* \right) \frac{\partial f_S}{\partial t},
\]

(9)

where \(D_i\) is the solute diffusion coefficient of component \(i\) (Cr, Mo, Nb, Al, Ti and C in the present case). In addition, the third term in Eq. (9) is the solute source term which is segregated at the solid-liquid interface and \(C_{L,i}^*\) and \(C_{S,i}^*\) are the solute concentrations of component \(i\) at the solid-liquid interface, respectively. Under equilibrium conditions, the two solute concentration terms at the solid-liquid interface are related as follows:

\[
C_{S,i}^* = k_i C_{L,i}^*,
\]

(10)

where \(k_i\) is the equilibrium partition coefficient of element \(i\). However, for the rapid solidification rate considered in the present study, the non-equilibrium relation between the element concentrations in the liquid and solid phases is given as:

\[
C_{S,i}^* = k_i \frac{a_0 R_s}{D_i} C_{L,i}^*,
\]

(11)

where \(a_0\) is the length of atomic dimensions of the components and \(R_s\) is the solidification rate, as computed by the CA model. From Eq. (11), when the solidification speed is slow, \(R_s \to 0\) and \(C_{S,i}^* = k_i C_{L,i}^*\). However, when the solidification speed is very fast, \(R_s \to \infty\) and \(C_{S,i}^* = C_{L,i}^*\). In other words, no element segregation occurs since the elements are captured by the solid phase before they can diffuse to the liquid phase.

Algorithm of CA model

The simulation domain is divided into cells with a uniform size of 0.2 \(\mu m \times 0.2 \mu m\). Moreover, the temperature at each point in the domain is considered to be both time- and position-dependent and is calculated as:

\[
T(z, t) = T_L - \hat{T} t + G_z Z,
\]

(12)

where \(\hat{T}\) is the cooling rate and \(G_z\) is the temperature gradient. The cooling rate can be estimated as

\[
\hat{T} \approx \frac{2\pi k_T e}{\epsilon} (T - T_0),
\]

(13)

where \(k_T\) is the thermal conductivity of IN713LC, \(e\) is the energy density, and \(T_0\) is the reference temperature. In addition, the temperature gradient can be calculated as the following space deviation:

\[
G_z = \frac{\hat{T} z}{2a}.
\]

(14)

It is noted that \(\hat{T}\) and \(G_z\) both depend on the energy density. In the CA model, each cell is characterized by two variables, namely the state (i.e., solid or liquid) and the crystallography orientation. In initializing the model, the system is considered to be at the liquidous temperature, and hence all of the cell states are set as 0. However, as the temperature of each cell falls below the liquidus temperature, the cell state is changed from 0 to 1; indicating that the cell resides in the solid state. Notably, the cell state may change from 0 or 1 under the following two conditions:

- Nucleation occurs in the cell
The cell is captured by a solid cell. The density of new nuclei formed during the solidification process is calculated by Eq. (1) for simulation domain at boundaries or inside the melt. The nucleation densities are then multiplied by the total numbers of cells with a state of 1 in order to determine the number of new nucleation sites. That is,

\[
N_s = \frac{n_{s_{\text{max}}}}{\sqrt{2\pi}} \exp \left[ -\frac{(\Delta T - \Delta T_{s_{\text{max}}})^2}{2\Delta T_{s_{\text{max}}}^2} \right] N_{as}, \quad (15-a)
\]

\[
N_v = \frac{n_{v_{\text{max}}}}{\sqrt{2\pi}} \exp \left[ -\frac{(\Delta T - \Delta T_{v_{\text{max}}})^2}{2\Delta T_{v_{\text{max}}}^2} \right] N_{av}, \quad (15-b)
\]

where \(N_{as}\) is the total number of cells with state 1 at the system boundary, and \(N_{av}\) is the total number of cells with state 1 inside the bulk liquid. Furthermore, \(\Delta T_{v_{\text{max}}}^0\) is considered to be 5 °C lower than \(\Delta T_{s_{\text{max}}}^0\). In performing the simulations, a random number, \(r_i\), with a value between 0 and 1 is assigned to each CA cell and nucleation is assumed to occur if the following condition is then satisfied:

\[
r_i < \frac{N_s}{N_{as}}. \quad (16)
\]

Following nucleation, a second random process is used to generate and assign a crystallography orientation number, \(q = [0 - 64]\), to the corresponding cell. Note that a large number of possible crystallography orientations is deliberately considered here in order to avoid impinging grains with the same orientation during the solidification process. In addition, the growth length of cell \(i\) with regard to its neighbor \(j\) at time \(t\) is calculated as

\[
L_i^j(t) = \int_0^t v_{\text{growth}} \, dt. \quad (17)
\]

Liquid neighbor capturing is also regarded as a random process. In particular, neighbor \(j\) is captured by cell \(i\) if the following condition is satisfied:

\[
r_i < \frac{L_i^j(t)}{\Delta x |\cos \theta_i + |\sin \theta_i| |}, \quad (18)
\]

where \(\Delta x\) is the cell length and \(\theta\) is the crystallographic orientation, which is calculated from the state of cell \(i\) as follows:

\[
\theta_i = \frac{\pi}{180} \left( 45 - 90 \frac{q_i}{q} \right). \quad (19)
\]

Equations 18 and 19 show that the grain growth speed is a function of the grain orientation. In particular, the grain orientation affects the capturing rate and then drives the grain growth phenomena. After cell \(i\) captures cell \(j\), the orientation of cell \(j\) is assigned to that of cell \(i\). The solid fraction of cells, \(f_s\), is then calculated at each time step using the following equation:

\[
f_s(t) = \frac{L_i^j(t)}{\Delta x}. \quad (20)
\]

Considering the need for convergence and numerical stability, the time step in the CA simulation process is limited as follows:

\[
dt \leq \frac{1}{5} \left( \frac{\Delta x}{df_s(t)/dt} \right)^{-\frac{1}{2}}. \quad (21)
\]
Materials and Methods

Laser processing

The IN713LC powder used in the present study had a spherical shape with a diameter of 15 - 60 μm. The printing trials were performed on a Tongtai AM-250 SLM machine equipped with a Nd-YAG laser with a wavelength of 1064 nm, a focal beam diameter of 100 μm, and a Gaussian irradiation profile. The laser provided a maximum power of 500 W and a maximum scanning speed of 2000 mm/s. The trials used a zig-zag scanning model with a 67° rotation of the scanning direction between consecutive layers. To avoid oxidation during the SLM process, the experiments were performed in a nitrogen atmosphere with an oxygen content of less than 1000 ppm. Based on a series of preliminary experiments, the scanning process was performed using three different laser energy densities of 360, 280 and 210 J/m³, respectively. For each energy density, three cubic specimens were printed with dimensions of 10 × 10 × 10 mm³.

Characterization

Following the SLM process, the samples were cut from the base plate with a wire and mounted on epoxy resin. The mounted samples were ground progressively with SiC sandpaper from a grit size of P240 to a final grit size of P3000, and were then further polished with a diamond suspension to a final particle size of around 0.3 μm. The polished samples were cleaned in ethanol solution and then dried. The microstructures and element compositions of the samples were examined using a scanning electron microscope (SEM, ZEISS Supra 55) equipped with energy dispersive X-ray spectroscopy (EDS) and electron backscattered diffraction (EBSD). In addition, the crack density was calculated using ImageJ software.

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**Author contributions statement**

Amir Reza Ansari Dezfoli has written the microstructure code. M. Nohsin Raza did experimental measurements. Yu-Lung Lo provided insight for the result discussion and edited the manuscript.
**Competing financial interests**

The authors declare no competing financial interests.