Riemannian optimal control

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Abstract

The aim of this paper is to adapt the general multitime maximum principle to a Riemannian setting. More precisely, we intend to study geometric optimal control problems constrained by the metric compatibility evolution PDE system; the evolution ("multitime") variables are the local coordinates on a Riemannian manifold, the state variable is a Riemannian structure and the control is a linear connection compatible to the Riemannian metric. We apply the obtained results in order to solve two flow-type optimal control problems on Riemannian setting: firstly, we maximize the total divergence of a fixed vector field; secondly, we optimize the total Laplacian (the gradient flux) of a fixed differentiable function. Each time, the result is a bang-bang-type optimal linear connection. Moreover, we emphasize the possibility of choosing at least two soliton-type optimal (semi-) Riemannian structures. Finally, these theoretical examples help us to conclude about the geometric optimal shape of pipes, induced by the direction of the flow passing through them.

Keywords: multitime maximum principle, Riemannian optimal control, shape optimization, gradient flow, total divergence, total Laplacian, bang-bang-type optimal solution, soliton-type metric.

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1 Adjoint PDE systems in Riemannian geometry

A connection on a manifold is a type of differentiation that acts on vector fields, differential forms and tensor products of these objects (see [2], [6], [11]). Its importance lies in the fact that given a piecewise continuous curve connecting two points on the manifold, the connection defines a linear isomorphism between the respective tangent spaces at these points. Another fundamental concept in the study of differential geometry is that of a Riemannian metric. It is well known that a Riemannian metric uniquely determines
a Levi-Civita connection: a symmetric connection for which the Riemannian metric is parallel. Since we may define linear connections unattached to Riemannian metrics, it is natural to ask, for a symmetric connection, if there exists a parallel Riemannian metric, that is, whether the connection is a Levi-Civita one. More generally, a connection on a manifold $M$, symmetric or not, is said to be metric if admits a parallel Riemannian metric defined on $M$. Then, the equations describing the metric property of symmetric linear connections are called metric compatibility evolution PDE system.

Let $M$ be an $n$-dimensional differentiable manifold with local coordinates $(x^1, ..., x^n)$. As we have mentioned above, fixing a Riemannian structure $g$ on $M$ ensures us about the existence of a symmetric linear connection satisfying the metric compatibility PDE system

$$\nabla_{\partial_i} g = 0, \forall i = 1, ..., n.$$ 

Let us change the geometric point of view in those of deformation theory. In this sense, let us discuss about the Riemannian metric controlled by a connection. For that we consider the controlled evolution law (linear PDE system)

$$(PDE) \quad \frac{\partial g_{ij}}{\partial x^k}(x) = g_{ps}(x) \left[ \delta^p_i \Gamma^s_{jk}(x) + \delta^p_j \Gamma^s_{ik}(x) \right], \quad i, j, k = 1, ..., n,$$

together with the initial condition

$$(x_0) \quad g_{ij}(x_0) = \eta_{ij},$$

where the piecewise metric tensor $g = (g_{ij})$ denotes a symmetric state tensor, $x = (x^1, ..., x^n)$ is the multitime variable (see [1], [3], [12]-[21]), and $\Gamma = (\Gamma^k_{ij})$ denotes the symmetric control linear connection.

The PDE system has solutions if and only if the complete integrability conditions

$$(CIC) \quad \frac{\partial}{\partial x^l} \left\{ g_{ps} \left[ \delta^p_i \Gamma^s_{jk} + \delta^p_j \Gamma^s_{ik} \right] \right\} = \frac{\partial}{\partial x^k} \left\{ g_{ps} \left[ \delta^p_i \Gamma^s_{jl} + \delta^p_j \Gamma^s_{il} \right] \right\}, \quad \forall i, j, k, l = 1, ..., n$$

are satisfied. Explicitly, this means $R_{ijkl} + R_{jikl} = 0$, where $R_{ijkl}$ denotes the Riemann curvature tensor field corresponding to the solution $(g, \Gamma)$.

We consider the set of admissible controls

$$\mathcal{U} = \{ \Gamma : M \to R^{n^3} | \Gamma^k_{ij} = \Gamma^k_{ji} \}.$$ 

Since the PDE system is linear, it coincides with its infinitesimal deformation, around a solution $g_{ij}(x)$. This PDE is also auto-adjoint since $v p_{\eta} - p v_{\eta} = 0$, for any two solutions $v(x, t)$ and $p(x, t)$. If it is used as adjoint equation, then a solution $p(x, t)$ is called the costate function.
The foregoing PDE systems determine the *multitime adjoint PDEs*

\[
(ADJ) \quad \frac{\partial \lambda_{ijk}}{\partial x^k}(x) = -\lambda_{psk}(x)[\delta^k_i \Gamma^{jk}_{sk}(x) + \delta^k_s \Gamma^{jk}_{pk}(x)],
\]

whose solution \( \lambda = (\lambda_{ijk}) \), called the *costate tensor*, is not necessary symmetric. The systems (PDE) and (ADJ) are adjoint (dual) in the sense of zero divergence of the vector field \( Q(x) = (Q^k(x) = y_{ij}(x)\lambda_{ijk}(x)) \), where \( y(x) = (y_{ij}(x)) \) denotes an infinitesimal deformation around a solution \( g_{ij}(x) \).

The symmetry of the state and control variables suggests us to consider a *symmetrized adjoint PDE system*, corresponding to the symmetric costate variables \( p_{ijk} = \lambda_{ijk} + \lambda_{jik} \). By computation, \( p \) is solution for

\[
(ADJ^s) \quad \frac{\partial p_{ijk}}{\partial x^k}(x) = -p_{rsk}(x)[\delta^k_i \Gamma^{jk}_{sk}(x) + \delta^k_s \Gamma^{jk}_{pk}(x)].
\]

Again, (PDE) and (ADJ^s) are dual in the sense of zero divergence of the vector field \( S(x) = (S^k(x) = y_{ij}(x)p_{ijk}(x)) \).

Moreover, introducing the Hamiltonian \( H(x, g, \Gamma, p) = p_{ijk}g_{st} \Gamma^{st}_{jk} \), the evolution systems (PDE) and (ADJ^s) become

\[
(PDE) \quad \frac{\partial g_{ij}}{\partial x^k}(x) = \left[ \frac{\partial H}{\partial p_{ijk}} + \frac{\partial H}{\partial p_{jik}} \right](x, g(x), \Gamma(x), p(x)),
\]

respectively

\[
(ADJ^s) \quad \frac{\partial p_{ijk}}{\partial x^k}(x) = -\left[ \frac{\partial H}{\partial g_{ij}} + \frac{\partial H}{\partial g_{ji}} \right](x, g(x), \Gamma(x), p(x)).
\]

# 2 Geometric optimal control with metric evolution-type constraints

## 2.1 Multitime maximum principle

In this section we consider a general control problem, with functional defined as multiple integral and evolution described by the Riemannian metric compatibility PDE system. Its solution is based on a general multitime maximum principle analyzed, under different aspects in [3] and [12]-[21], and generalizing the classical approach on single-time Pontriaguine maximum principle (see [4],[7]-[10],[22]). More precisely, in [21] we have proved, using needle-shaped control variations, that, given the hyper-parallelepiped \( \Omega_{0t} \) in \( \mathbb{R}^m \)
having 0 and $t_0$ as diagonal points, the solutions for the multitime optimal control problem with running cost and integral cost on boundary,

$$\max_{u(\cdot)} \left( \mathcal{J}[u(\cdot)] = \int_{\Omega_{t_0}} X(t, x(t), u(t)) \, dt + \int_{\partial \Omega_{t_0}} \chi(t, x(t)) \, d\sigma \right)$$

subject to 

$$\frac{\partial x^i}{\partial t}(t) = X^i_\alpha(t, x(t), u(t)), \ i = 1, \ldots, n, \ \alpha = 1, \ldots, m,$$

$$u(t) \in R^k, \ x(t) \in R^n, \ t \in \Omega_{t_0} \subset R^m, \ x(0) = x_0,$$

with the corresponding control Hamiltonian

$$H(t, x, p, u) = X(t, x, u) + p^\alpha X^i_\alpha(t, x, u),$$

satisfy the multitime maximum principle described in the next fundamental outcome.

**Theorem (Multitime maximum principle).** Suppose $u^*(\cdot)$ is an optimal solution of the control problem and $x^*(\cdot)$ is the corresponding optimal state. Then there exists a costate tensor $p^*(\cdot) : \Omega_{t_0} \rightarrow R^{mn}$ such that

$$\frac{\partial x^i}{\partial t}(t) = \frac{\partial H}{\partial p^\alpha}(t, x^*(t), p^*(t), u^*(t)),$$

$$\frac{\partial p^\alpha}{\partial t}(t) = -\frac{\partial H}{\partial x^i}(t, x^*(t), p^*(t), u^*(t)),$$

$$H(t, x^*(t), p^*(t), u^*(t)) = \max_{u(\cdot)} H(t, x^*(t), p^*(t), u(t)), \ \forall t \in \Omega_{t_0}$$

and

$$n_\alpha p^\alpha |_{\partial \Omega_{t_0}} = \frac{\partial \chi}{\partial x^i} |_{\partial \Omega_{t_0}}.$$
where \( dx \) denotes the differential \( n \)-form \( dx^1 \wedge ... \wedge dx^n \) (the Euclidean volume element on \( \Omega_{x_0,x_1} \)) and \( d\sigma \) is the Euclidean volume element on the boundary.

The multitime optimal control problem consists in finding

\[
(J) \quad \max_{\Gamma} \left( J[\Gamma] = \int_{\Omega_{x_0,x_1}} X(x, g(x), \Gamma(x)) \, dx + \int_{\partial \Omega_{x_0,x_1}} \chi(x, g(x)) \, d\sigma \right)
\]

subject to the evolution system

\[
(PDE) \quad \frac{\partial g_{ij}}{\partial x^k}(x) = g_{ps}(x) \left[ \delta_p^i \Gamma^s_{jk}(x) + \delta_p^j \Gamma^s_{ik}(x) \right], \quad i, j, k = 1, ..., n
\]

and the initial condition

\[
(x_0) \quad g(x_0) = \eta.
\]

Since the main ingredients of this Riemannian optimal control problem (the state variables, the control variables and evolution constraints \((PDE)\)) are symmetric, we shall derive an adapted multitime maximum principle, based on symmetric costate variables. For this, we introduce the symmetric Lagrange multipliers \( p^{ijk} = p^{jik} \) and the reduced control Hamiltonian

\[
(H) \quad H(x, g, \Gamma, p) = X(x, g, \Gamma) + g_{is} \Gamma^s_{jk} p^{ijk}.
\]

**Corollary 1 (Riemannian maximum principle).** Suppose the symmetric connection \( \Gamma^*(\cdot) \) is an optimal solution for \((PDE), (J), (x_0)\) and that \( g^*(\cdot) \) is the corresponding optimal Riemannian structure. Then there exists a symmetric dual tensor \( p^* = (p^{*ijk}) : \Omega_{x_0,x_1} \to \mathbb{R}^n \) such that

\[
(PDE) \quad \frac{\partial g^*_{ij}}{\partial x^k}(x) = \left[ \frac{\partial H}{\partial p^{*jk}} + \frac{\partial H}{\partial g^{*ik}} \right](x, g^*(x), \Gamma^*(x), p^*(x)),
\]

\[
(ADJ^*) \quad \frac{\partial p^{*ijk}}{\partial x^k}(x) = - \left[ \frac{\partial H}{\partial g^*_{ij}} + \frac{\partial H}{\partial g^{*ji}} \right](x, g^*(x), \Gamma^*(x), p^*(x))
\]

and

\[
(Opt) \quad H(x, g^*(x), \Gamma^*(x), p^*(x)) = \max_{\Gamma(\cdot)} H(x, g^*(x), \Gamma(x), p^*(x)), \quad \forall x \in \Omega_{x_0,x_1}.
\]

Finally, the boundary conditions

\[
(\partial_{\Omega_{x_0,x_1}}) \quad n_k p^{*ijk} |_{\partial_{\Omega_{x_0}}} = \left[ \frac{\partial \chi}{\partial g_{ij}} + \frac{\partial \chi}{\partial g^{*ji}} \right]_{\partial_{\Omega_{x_0,x_1}}}
\]
are satisfied, where \( n \) denotes the covector corresponding to the unit normal vector on \( \partial \Omega_{x_0 x_1} \).

**Proof.** We denote by \( \mathcal{H} \) the standard control Hamiltonian corresponding to the Riemannian optimal control problem \( ((J), (PDE), (x_0)) \), that is

\[
\mathcal{H}(x, g, \Gamma, \lambda) = X(x, g, \Gamma) + \lambda^{ij} \left[ g_{is} \Gamma_{jk}^s(x, g, \Gamma) + g_{js} \Gamma_{ik}^s(x, g, \Gamma) \right]
\]

\[
= X(x, g, \Gamma) + g_{is} \Gamma_{jk}^s \left[ \lambda^{ij} + \lambda^{jk} \right] = H(x, g, \Gamma, \lambda^{ij} + \lambda^{jk}).
\]

Let us define the symmetric costate tensor \( p^{ijk} = \lambda^{ijk} + \lambda^{jik} \). Writing the multitime maximum principle with standard Hamiltonian, and using the definition of \( p \), we obtain

\[
\frac{\partial g^{ij}_{\ast}}{\partial x^k} = \frac{\partial \mathcal{H}}{\partial \lambda^{ij}} = \frac{\partial H}{\partial p^{ijk}} + \frac{\partial H}{\partial p^{jik}};
\]

\[
\frac{\partial p^{ijk}_{\ast}}{\partial x^k} = \frac{\partial \lambda^{ijk}_{\ast}}{\partial x^k} + \frac{\partial \lambda^{jik}_{\ast}}{\partial x^k} = - \left[ \frac{\partial H}{\partial g_{ij}} + \frac{\partial H}{\partial g_{ji}} \right] = - \left[ \frac{\partial H}{\partial g_{ij}} + \frac{\partial H}{\partial g_{ji}} \right];
\]

\[
H(x, g^{\ast}, \Gamma^{\ast}, p^{\ast}) = \mathcal{H}(x, g^{\ast}, \Gamma^{\ast}) = \max_p \mathcal{H}(x, g^{\ast}, \Gamma^{\ast}, p^{\ast});
\]

\[
n_k p^{ijk}_{\ast}|_{\partial \Omega_{x_0 x_1}} = n_k \left[ \lambda^{ijk}_{\ast} + \lambda^{jik}_{\ast} \right]_{\partial \Omega_{x_0 x_1}} = \left[ \frac{\partial \chi}{\partial g_{ij}} + \frac{\partial \chi}{\partial g_{ji}} \right]_{\partial \Omega_{x_0 x_1}}.
\]

**Remark.** By replacing the metric compatibility evolution \( (PDE) \) with the PDE system corresponding to the dual tensor \( g^{-1} \):

\[
(PDE') \quad \frac{\partial g^{ij}_{\ast}}{\partial x^k}(x) = -g^{ps}(x) \left[ \delta^i_p \Gamma^{j}_{sk}(x) + \delta^j_p \Gamma^{i}_{sk}(x) \right], \quad i, j, k = 1, \ldots, n,
\]

with initial condition

\[
(x_0') \quad g^{ij}(x_0) = \eta^{ij}
\]

and using the dual Hamiltonian

\[
(H') \quad H'(x, g^{-1}, \Gamma, p) = X(x, g, \Gamma) - g^{is} \Gamma^{jk}_{sk} p^{kj}_{ij},
\]

we can rephrase the Riemannian multitime maximum principle as it follows.

**Corollary 2 (Riemannian dual maximum principle).** Suppose the symmetric connection \( \Gamma^{\ast}(\cdot) \) is an optimal solution for \( ((PDE'), (J), (x_0')) \) and that \( g^{s-1} \) is the corresponding optimal state. Then there exists a symmetric dual tensor \( p^{\ast} = (p^{kj}_{ij}) : \Omega_{x_0 x_1} \to \mathbb{R}^{n^3} \) such that

\[
(PDE') \quad \frac{\partial g^{s-1}_{ij}}{\partial x^k}(x) = \left[ \frac{\partial H'}{\partial p^{kj}_{ij}} + \frac{\partial H'}{\partial p^{kj}_{ij}} \right](x, g^{s-1}(x), \Gamma^{\ast}(x), p^{\ast}(x)),
\]
\[
(ADJ^s) \quad \frac{\partial p^s_{ij}}{\partial x^k}(x) = - \left[ \frac{\partial H'}{\partial g^{ij}} + \frac{\partial H'}{\partial g^{ji}} \right] (x, g^{s-1}(x), \Gamma^s(x), p^s(x))
\]

and

\[
(Opt) \quad H'(x, g^{s-1}(x), \Gamma^s(x), p^s(x)) = \max_{\Gamma(\cdot)} H'(x, g^{s-1}(x), \Gamma(x), p^s(x)), \quad \forall x \in \Omega_{x_0 x_1}.
\]

Finally, the boundary conditions

\[
(\partial \Omega_{x_0 x_1}) \quad n_k p^s_{ij} |_{\partial \Omega_{x_0 x_1}} = \left[ \frac{\partial \chi}{\partial g^{ij}} + \frac{\partial \chi}{\partial g^{ji}} \right] \frac{\partial}{\partial x^k} |_{\partial \Omega_{x_0 x_1}}
\]

are satisfied.

3 Flux-type optimal control problems

Throughout this section, the basic geometric ingredients have the same significance as above; that is, \((M, g)\) is a Riemannian manifold with local coordinates \(x = (x^1, ..., x^n)\) and \(\Omega_{x_0 x_1}\) denotes a subset of \(M\) diffeomorphic with a hyper-parallelepiped in \(\mathbb{R}^n\). The main goal of the section consists in analyzing two flux-type Riemannian optimal control problems, resulting in bang-bang-type optimal solutions. The key idea is to take \(x = (x^1, ..., x^n)\) like an evolution (deformation) parameter.

3.1 Optimization of total divergence

In this subsection, \(X\) is a fixed vector field on \(\Omega_{x_0 x_1}\). The optimal control problem we are looking to solve consists in finding the control connection

\[
(\Gamma^k_{ij}) \in \mathcal{U} = \{ \Gamma : \Omega_{x_0 x_1} \to [-1, 1]^{n^3} | \Gamma^k_{ij} = \Gamma^k_{ji} \}
\]

that maximize the total divergence of \(X\). More precisely, we try to find the optimal linear connection and the optimal Riemannian structure that maximize the Bolza-type functional

\[
(J) \quad J[\Gamma] = \int_{\Omega_{x_0 x_1}} \text{Div } X \, dv = \int_{\Omega_{x_0 x_1}} \text{Div } X \sqrt{g} \, dx,
\]

where \(g = \det(g_{ij})\), subject to the dual metric compatibility evolution PDE system

\[
(PDE') \quad \frac{\partial g^{ij}}{\partial x^k}(x) = -g^{ps}(x) \left[ \delta_p^i \delta_s^j (x) + \delta_p^j \delta_s^i (x) \right], \quad i, j, k = 1, ..., n,
\]
with initial condition
\[(x_0')\quad g^{ij}(x_0) = \eta^{ij}.\]

**Remark.** By applying the *Divergence Theorem* in Riemannian setting, we may rewrite the functional \(J[\Gamma]\) as
\[
J[\Gamma] = \int_{\partial \Omega} g(X, N^g) d^g\sigma,
\]
or, using local coordinates,
\[
J[\Gamma] = \int_{\partial \Omega} X^i n_i \sqrt{g} d\sigma.
\]
In the expressions above, if \(N^g = (N^i)\) is the outpointing normal vector field on the boundary, with respect to the metric \(g\), then \(n = (n_i = g_{ij}N^j)\), denotes the Euclidean normal covector along the boundary. This allows us to identify the running cost, respectively the boundary cost associated to this optimal control problem:
\[
X(x, g^{-1}, \Gamma) = 0; \quad \chi(x, g^{-1}) = X^i \sqrt{g} n_i,
\]
giving the corresponding control Hamiltonian
\[
H'(x, g^{-1}, \Gamma, p) = -g^{ij}\Gamma_{sk}^l p_{ij}^k.
\]
Since this Hamiltonian is linear with respect to the control components \(\Gamma_{ij}^k\), we have no interior optimal control \(\Gamma_{ij}^k\); for optimum, the control must be at a vertex of \([-1, 1]^n\) (see linear optimization, simplex method).

Writing the adjoint PDE system
\[
\frac{\partial p_{ij}^k}{\partial x^k} = p_{is}^k \left[ \delta_i^l \Gamma_{jk}^s + \delta_j^l \Gamma_{ik}^s \right],
\]
we obtain the immediate solution \(p_{ij}^k = C_{ij}^k g_{ij}^s\), with \(C = (C^k) : \Omega_{x_0x_1} \to \mathbb{R}^n\), \(\frac{\partial C^k}{\partial x^k} = 0\). Then
\[
H'(x, g^{-1}, \Gamma, p^*) = -C_{ij}^k \Gamma_{ks}^e.
\]
Therefore, the optimal control maximizing the total divergence is a linear connection having the bang-bang-type components
\[
\Gamma_{ij}^* = \begin{cases} 
\delta_{ij}^l & \text{if } k = i, \epsilon^j \neq 0 \\
\delta_{ij}^l & \text{if } k = j, \epsilon^i \neq 0 \\
\text{arbitrary,} & \text{otherwise}
\end{cases}
\]
where \( \epsilon_l = \text{sgn}(-C^l) \).

Moreover, the boundary constraints corresponding to these solutions are

\[
[n_k C^k g^*_i](x) = [n_k (\sqrt{g^* X^k}) g^*_i](x), \quad \forall x \in \partial \Omega_{x_0 x_1},
\]

that is

\[
n_k(x) \left( C^k - \sqrt{g^* X^k} \right)(x) = 0, \quad \forall x \in \partial \Omega_{x_0 x_1}
\]

and, together with the initial condition

\[
g^{-1}(x_0) = \eta
\]

may help us to determine the solenoidal vector field \( C \).

**Remark.** If we replace the above maximum-type problem with a minimizing one, we obtain similar solutions, with \( \epsilon_l = \text{sgn}(C^l), \ l = 1, \ldots, n \).

### 3.2 Optimization of total Laplacian

Let \( f : \Omega_{x_0 x_1} \to \mathbb{R} \) be a fixed differentiable function. The optimal control problem we are interested in consists in maximizing the functional

\[
(J) \quad J[\Gamma] = \int_{\Omega_{x_0 x_1}} \Delta g f dv = \int_{\partial \Omega_{x_0 x_1}} g^{ij} f_i n_j \sqrt{g} d\sigma,
\]

subject to the dual metric compatibility evolution PDEs system

\[
(PDE') \quad \frac{\partial g^{ij}}{\partial x^k}(x) = -g^{ps}(x) \left[ \delta^i_p \Gamma^j_{sk}(x) + \delta^j_p \Gamma^i_{sk}(x) \right], \quad i, j, k = 1, \ldots, n,
\]

with control restriction

\[
(\Gamma^k_{ij}) \in \mathcal{U} = \{ \Gamma : \Omega_{x_0 x_1} \to [-1, 1]^n \mid \Gamma^k_{ij} = \Gamma^k_{ji} \}
\]

and with initial condition

\[
(x'_0) \quad g^{ij}(x_0) = \eta^{ij}.
\]

The running cost, respectively the boundary cost associated to this optimal control problem are

\[
X(x, g^{-1}, \Gamma) = 0; \quad \chi(x, g^{-1}) = g^{ij} \sqrt{g} f_i n_j,
\]

where \( f_k(x) = \frac{\partial f}{\partial x^k} \) and \( g^{ij} = g^{ij}(x) \) denote the components of the inverse metric matrix. Again, the control Hamiltonian is

\[
H'(x, g^{-1}, \Gamma, p) = -g^{is} \Gamma^j_{sk} p^k_{ij}
\]
and, since it is linear with respect to the control components $\Gamma_{ij}^k$, we have no interior optimal control $\Gamma_{ij}^k$.

Writing the Riemannian maximum principle gives us the adjoint PDE system

$$\frac{\partial p_{ij}^k}{\partial x^k} = p_{is}^k \left[ \delta_i^j \Gamma_{jk}^s + \delta_j^i \Gamma_{ik}^s \right],$$

with same possible solution as in the previous section, that is $p_{ij}^k = C^k g_{ij}^*$, with $C = (C^k) : \Omega_{x_0 x_1} \rightarrow \mathbb{R}^n, \frac{\partial C^k}{\partial x^k} = 0$. Replacing within the control Hamiltonian, we obtain

$$H'(x, g^{*^{-1}}, \Gamma, p^*) = -C^a \Gamma_{as}.$$ 

Therefore, the optimal control maximizing the gradient flux is of bang-bang-type

$$\Gamma_{ij}^* = \begin{cases} \delta_{jl} \xi_l & \text{if } k = i, \xi \neq 0 \\ \delta_{il} \xi_l & \text{if } k = j, \xi \neq 0 \\ \text{arbitrary, otherwise,} \end{cases}$$

where $\xi_l = \text{sgn}(-C^l)$.

This time instead, the boundary constraints generating the solenoidal tensor field $C$ are

$$[n_k C^k g_{ij}^*](x) = [n_k g^{*kl} \left( f_{il} g_{lj}^* + f_{jl} g_{li}^* - f_{ll} g_{ij}^* \right) \sqrt{g}] (x), \quad \forall x \in \partial \Omega_{x_0 x_1}.$$ 

**Remark.** For both the foregoing problems, we may look for some particular solutions.

1. We may chose

$$\Gamma_{ij}^* = \delta_{ij}^k (\delta_{jl} \xi_l) + \delta_{ik}^j (\delta_{il} \xi_l) - \delta_{ij} \delta^{kp} (\delta_{pl} \xi_l),$$

that is $\Gamma^*$ is an Euclidean conformal linear connection (see [5]), i.e. $\Gamma^*$ is conformal with the Levi-Civita connection associated to the Euclidean metric $g_{ij}^0 = \delta_{ij}$. Then, the optimal Riemannian metric

$$g_{ij}^{*ij} = K \delta_{ij} e^{-2 \delta_{kl} \xi^k \xi^l}$$

is a soliton-type solution for the dual metric compatibility evolution (PDE)

$$\frac{\partial g_{ij}}{\partial x^k}(x) = -g^{*ps}(x) \left[ \delta_i^p \Gamma_{sj}^k(x) + \delta_j^p \Gamma_{is}^k(x) \right]$$

and, also, is a dual Riemannian structure (a dual Riemannian metric having $\Gamma^*$ as Levi-Civita connection).
2. We may consider
\[ \Gamma^k_{ij} = \epsilon^k \epsilon_i \epsilon_j, \]
where \( \epsilon_i = \delta_{ij} \epsilon^j \). Then, the \((PDE')\) system writes
\[ \frac{\partial g^{ij}}{\partial x^k} = \left( g^{is} \epsilon^j + g^{js} \epsilon^i \right) \epsilon_s \epsilon_k, \]
admitting the following soliton-type solution
\[ g^{ij} = \left[ \alpha e^{-2n \epsilon_k x^k} + \frac{\alpha_i + \alpha^j}{2} e^{-n \epsilon_k x^k} \right] \epsilon^i \epsilon^j \quad \text{(no summation),} \]
where \( \alpha, \alpha^i \) denote real constants, satisfying \( \sum_{i=1}^n \alpha^i = 0 \). The disadvantage of the latter solution is that it may not be a Riemannian structure, but only a symmetric \((2,0)\)-type tensor field or, in best case scenario (i.e. \( \epsilon^i \neq 0, \forall i = 1, ..., n \)), a semi-Riemannian structure.

4 The optimal geometry of pipes

This section is meant to emphasize the practical utility of the theoretical facts described above, by analyzing a classical problem in Hydraulics and Fluid Mechanics. Given a pipe, in the general sense (water pipe, gas pipe, blood vessel) containing a fluid flow, it is well known that the Divergence Theorem helps us to measure the flux of the fluid flow through pipe walls. Sometimes instead, for practical reasons, it is of major utility to identify the optimal shape of the pipe, allowing the minimum flux through walls. This is the problem analyzed in this section. More precisely, given the directionality of the fluid through the pipe, we decide about the best way to conceive the pipe (the optimal geometric shape), such that the flux of the fluid through pipe walls to be minimal.

Let \( D^1 \) denote the closed disc of radius one and let \( M = D^1 \times (0,1) \) be a differential manifold with boundary describing the interior and the boundary of a cylinder. We identify the pipe in the Euclidean space (in the sense of some diffeomorphism) with the manifold \( M \). Given a flow through the pipe, described by a vector field \( F = X \partial_x + Y \partial_y + Z \partial_z \) on \( M \), we shall find a Riemannian structure on \( M \), minimizing the flux of \( F \). For this, we consider the local map \( V = M - \{(x,y,z) \in M \mid y = 0, x \geq 0\} \). Using the cylindrical coordinates \((\rho, \theta, z)\), we may identify \( V \) with the parallelepiped \((0,1] \times (0,2\pi) \times (0,1) \) in \( \mathbb{R}^3 \). Moreover, we suppose that the expression of \( F \)
with respect to these new coordinates is \( F = R \frac{\partial}{\partial \rho} + T \frac{\partial}{\partial \theta} + \zeta \frac{\partial}{\partial z} \). Then,

\[
\begin{align*}
R(\rho, \theta, z) &= X(\rho \cos \theta, \rho \sin \theta, z) \cos \theta + Y(\rho \cos \theta, \rho \sin \theta, z) \sin \theta; \\
T(\rho, \theta, z) &= \rho \left[ -X(\rho \cos \theta, \rho \sin \theta, z) \sin \theta + Y(\rho \cos \theta, \rho \sin \theta, z) \cos \theta \right]; \\
\zeta(\rho, \theta, z) &= Z(\rho \cos \theta, \rho \sin \theta, z),
\end{align*}
\]

or, conversely,

\[
\begin{align*}
X(x, y, z) &= R(\sqrt{x^2 + y^2}, \arctan \frac{y}{x}, z) \frac{x}{\sqrt{x^2 + y^2}} - T(\sqrt{x^2 + y^2}, \arctan \frac{y}{x}, z)y; \\
Y(x, y, z) &= R(\sqrt{x^2 + y^2}, \arctan \frac{y}{x}, z) \frac{y}{\sqrt{x^2 + y^2}} + T(\sqrt{x^2 + y^2}, \arctan \frac{y}{x}, z)x; \\
Z(x, y, z) &= \zeta(\sqrt{x^2 + y^2}, \arctan \frac{y}{x}, z).
\end{align*}
\]

Applying the results derived in the previous section, we obtain the optimal Euclidean conformal structure

\[
g(\rho, \theta, z) = Ke^{2\text{sgn}(R(1, \theta, z))\rho} \left( d\rho^2 + d\theta^2 + dz^2 \right).
\]

Using the above relations between the components of \( F \) relative to the cylindrical and Cartesian coordinates we derive also the Cartesian expression of the optimal Riemannian structure:

\[
g = Ke^{2S} \left( \sqrt{\frac{x}{\sqrt{x^2 + y^2}} \sqrt{\frac{y}{\sqrt{x^2 + y^2}}}} \right) \sqrt{x^2 + y^2} \begin{pmatrix}
1 & 0 & 0 \\
0 & x^2 + y^2 & 0 \\
0 & 0 & 1
\end{pmatrix},
\]

where \( S = \text{sgn}(\langle N, F \rangle) \) on the boundary, \( \langle ., . \rangle \) denotes the canonical inner product on \( \mathbb{R}^3 \) and \( N(x, y, z) = x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} \) is the normal vector field along the boundary of \( M \).

In conclusion, the direction of the flow \( F \) as vector field in \( \mathbb{R}^3 \) is directly involved in the geometric configuration of the pipe. More precisely, if \( F \) is pointed outward, then the diameter of the pipe increases; conversely, if \( F \) is pointed inward, the diameter decreases. Therefore, the optimal shape for the pipe walls is the one tangent, at each point, to the flow \( F \).

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