Detection of Solar Rotational Variability in the Large Yield RAdiometer (LYRA) 190 – 222 nm Spectral Band

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Abstract We analyze the variability of the spectral solar irradiance during the period from 7 January 2010 until 20 January 2010 as measured by the Herzberg channel (190 – 222 nm) of the Large Yield RAdiometer (LYRA) onboard PROBA2. In this period of time, observations by the LYRA nominal unit experienced degradation and the signal produced by the Herzberg channel frequently jumped from one level to another. Both factors significantly complicate the analysis. We present the algorithm that allowed us to extract the solar variability from the LYRA data and compare the results with SORCE/SOLSTICE measurements and with modeling based on the Code for the Solar Irradiance (COSI).
1. Introduction

After the start of regular space-born measurements in 1978, it became clear that the solar irradiance is variable on different time scales. Since then significant progress was achieved both in measuring and in modeling of the solar irradiance variability (Fröhlich, 2005; Krivova and Solanki, 2008). At the same time, new missions devoted to monitoring of the spectral solar irradiance (SSI) continue to bring significant surprises (Harder et al., 2009). So our understanding of the mechanisms of the solar irradiance variability depends on observations by new instruments. The importance of studying the solar irradiance variability is further emphasized by its direct impact on the Earth’s climate (Haigh et al., 2010; Gray et al., 2010).

In this article we analyze the spectral solar irradiance data from the recent European mission (PROBA2) (launched on 2 November, 2009). Large Yield RAdiometer (LYRA) is a solar radiometer onboard PROBA2, which is a technologically oriented ESA micromission, and is observing the solar irradiance in two UV and two EUV spectral channels (Hochedez et al., 2006; Benmoussa et al., 2009; Dominique et al., 2012b). The passbands of the UV channels were selected for their relevance for the ozone concentration.

Up to now the LYRA data were used only for the analysis of variations shorter than one day. Van Doorsselaere et al. (2012) studied flares observed by LYRA. Dominique et al. (2012a) and Shapiro et al. (2011c) analyzed the light curves during occultations and eclipses, respectively. In this article we analyze the variations of the solar irradiance during a period of approximately two weeks, which corresponds to the transit of active regions across the solar disk and is representative for studying the solar rotational cycle (Fligge, Solanki, and Unruh, 2000). We compare our results with SOLar-STellar Irradiance Comparison Experiment (SORCE/SOLSTICE: McClintock, Rottman, and Woods, 2005). We also use Code for the Solar Irradiance (COSI: Shapiro et al., 2010) as a tool for modeling the variability of the irradiance, assuming that it is determined by the evolution of the solar surface magnetic field. The theoretical results are compared with the LYRA measurements. Our analysis is restricted to the Herzberg channel of LYRA as irradiance in the Herzberg continuum (190 – 222 nm) presents challenge for the solar radiative transfer models (Shapiro et al., 2011c) and is especially important for climate modeling (Brasseur et al., 1987; Rozanov et al., 2006; Shapiro et al., 2011b).

The analysis of the solar variability in this channel is significantly hampered by the severe degradation, that LYRA started to experience immediately after the covers had been opened. The degradation led to a significant loss of sensitivity to the solar signal already after a month since first light (6 January 2010). This is why, for the current analysis, we used early data, although the instrument was still in its commissioning phase. At that time, the spacecraft was slightly off-pointing, due to a problem in its onboard software that was fixed at the end of January 2010. Another significant problem was therefore pointing fluctuations which in combination with inhomogeneous sensitivity of the diamond detectors resulted in significant fluctuations of the data. We note that the amplitude of the signal introduced by the jitter of PROBA2 exceeds the natural solar variability. The data were also corrupted by the angular rotation of PROBA2 (four times per orbit), occultations, electronic perturbations, and sudden jumps (which are specific for the Herzberg channel and are not yet fully understood). The latter problems could usually be easily identified and corrected, so they did not present such important difficulties as degradation and pointing fluctuations. The detailed discussion of the different perturbations in the LYRA data are given by Dominique et al. (2012b).

We analyze the LYRA data for the period from 7 January 2010 until 20 January 2010. During this period one sunspot group surrounded by a plage region made a full transit across
Figure 1  SOHO/MDI continuum images for 8 January (a), 11 January (b), 14 January (c), and 17 January (d), 2010. For better clarity the contrasts between sunspot, plage, and quiet Sun were artificially increased.

The sunspot group surrounded by plage appeared on the solar disk on 7 January 2010 and disappeared on 20 January 2010. The transit is shown on Figure 1. One can expect that the presence of active regions on the solar disk will modify the solar irradiance (Fligge, Solanki, and Unruh, 2000). In this section we present the method that was used to extract these modulations from the LYRA data.

The level3 calibrated data from the Herzberg channel of LYRA for January 2010 are plotted in Figure 2. These data are available for the community (Dominique et al., 2012b) and corrected for the temperature effects, degradation, dark current, and one-minute averaged. The LYRA samples in Figure 2 can be sorted in two groups: the bottom line corresponds to data acquired in occultations by the Earth (dark current), the upper series constitute the actual Herzberg time series, but is perturbed by jumps induced by the electronics (appearing when the FPGA was reloaded) and by pointing fluctuations. The selection of valid data was therefore not trivial, and we had to process them with a special care.

The part of the data corresponding to periods of significant pointing fluctuations or occultations had to be excluded from the analysis. Therefore, in the first step we chose a reliable range for the irradiance and excluded the outliers. For the period under consideration (January 2010) the lower and upper levels were chosen to be 0.65 Wm$^{-2}$ and 0.9 Wm$^{-2}$, respectively. This procedure is illustrated in Figures 3 and 4. One can see that even within the
Figure 2  Irradiance measured by the Herzberg channel of LYRA for the 6 January – 24 January 2010 period. Plotted are the level3 calibrated data.

selected trustable range the irradiance level is highly unstable and undergoes a few jumps per day. For some of the days, several distinctive “branches” can be clearly seen. The irradiance level is stable within each of the branches but constantly jumps between them. Such behavior can be attributed to pointing fluctuations and rotation of PROBA2, as well as to some electronic disturbances.

The upward slope in the zero level of the irradiance is due to a constant which was added to remove the degradation for the production of the level3 data. The additive correction of the degradation is justified for the analysis of the flares but is not suitable for the analysis of the rotational cycle, where a multiplicative correction is preferable (see Dominique et al., 2012b; Shapiro et al., 2011c). Therefore we had to subtract this constant from the level3 data. To correct for the degradation, we calculated the change of the irradiance at the Herzberg channel between 7 January and 18 January as measured by SORCE/SOLSTICE. The SOLSTICE data were convolved with the profile of the LYRA Herzberg channel (Benmoussa et al., 2009). Then, following the approach of Dominique et al. (2012b) for the second half year of LYRA observations, we tested two different corrections for the degradation, one of the type $\frac{1}{a+bt_{\text{time}}}$, and another of the type $\exp(a - bt_{\text{time}})$. A spline correction was not used as we are interested only in a relatively short period of time. In contrast to Dominique et al. (2012b), we applied the corrections not as an additive but as a multiplicative factor. In both cases we chose the coefficients $a$ and $b$ so that the change of the irradiance between 7 January and 18 January was the same in LYRA and SOLSTICE data. No substantial difference was seen between the data corrected with linear and exponential functions so we only show the data processed with a linear correction.

The segregation of the LYRA time series to different branches is further illustrated in Figure 5, which shows four selected time intervals with high temporal resolution. The data plotted in Figure 5(a) consists only of one branch which covers the 14.75 – 14.95 time interval. The data in the 14.95 – 15.05 interval were considered to be too noisy and excluded from the analysis. The data plotted in Figure 5(b) consist of two branches (both in the 15.44 – 15.60 interval), while the data in the 15.30 – 15.44 interval were excluded. The data plotted in Figure 5(c) consist of four branches (two branches prior to 15.99 and two after), while the data in Figure 5(d) consist of two branches covering the entire time interval.

The structure of the LYRA data significantly complicates the analysis as the amplitude of the jumps is sometimes larger than expected amplitude of the solar variability. The fundamental assumption of our analysis is that, although the variability of the LYRA data is
dominated by these jumps, every individual branch does have a signature of the solar variability.

To extract this signature, we binned the available data to four-orbit intervals (see Figures 6 and 7). Each interval starts when PROBA2 passes the Earth’s equatorial plane and lasts $4 \times 100 = 400$ minutes. In each of these intervals we took the first stretches of every branch (16.142 – 16.162 and 16.162 – 16.182 intervals on Figure 5(d)) and calculated the mean values of the irradiance in these stretches. Then we shifted one (or more if the interval contained more than two branches) of the branches enforcing the condition that these mean values should be equal to each other. The amplitude of the jumps did not change signifi-

Figure 3  Irradiance measured by the Herzberg channel of LYRA for the 6 January – 14 January 2010 period. The red solid lines denote the reliable range of the irradiance, which was used for the analysis.
Figure 4  Irradiance measured by the Herzberg channel of LYRA for the 14 January–22 January 2010 period. The red solid lines denotes the trustable range of the irradiance which was used for the analysis.

cantly within the four-orbit interval (which also confirms that jumps occur due to the fact that PROBA2 needs to rotate four times per orbit), so this allowed us to eliminate all of the jumps from the data. The procedure described above is basically equivalent to the correction of the flat field. In principle, one can expect that the mean values of the irradiance at the considered stretches could be different due to the solar variability. However the duration of almost all stretches was very small (Figure 5) so neglecting the variations of the solar irradiance between them does not lead to a significant error.

After the jumps in the data had been corrected, we applied a linear regression to all available data for each of the four-orbit intervals (the red lines in Figures 6 and 7). The slope
of the regression was used to calculate the change of the irradiance level during each of the intervals. In many cases the data did not cover the entire interval and contained significant gaps in the beginning or at the end of the interval. For simplicity we assumed that the level of the solar irradiance did not change during these gaps, so the change of the irradiance corresponds to the projection of the red lines in Figures 6 and 7 to the vertical axis. An alternative would be an extrapolation of the linear trend, however such extrapolation can significantly increase errors caused by the noise in the data.

Our final product is a single value of the irradiance per four-orbit interval. The differences between each of the consecutive values were calculated employing the algorithm described above. Let us note that the implementation of this algorithm violates the condition of the equal change of the irradiance between 7 January and 18 January as measured by LYRA and SOLSTICE and thus the correction for the degradation should be slightly readjusted. This readjustment had not yet been implemented in Figures 6 and 7, so almost all data in these figures show the upward trend.

The data in the 15.04–15.32 and 15.87–16.15 time intervals were too noisy even after the corrections. At the same time the irradiance significantly decreases during these intervals, which can therefore be just an artifact caused by noise. To take this into account we produced two different data sets of the LYRA data. The LYRA version1 data set was pro-
produced taking all data into account, while LYRA version2 data set was produced excluding the two aforementioned intervals from the analysis and assuming that the solar irradiance did not change during these intervals. The difference between these data sets indicates the accuracy of our analysis.

3. Modeling with COSI

In this section we calculate the synthetic profile of the spectral solar irradiance variability for the period analyzed in Section 2. We follow a well-developed approach (Foukal and Lean, 1988; Fligge, Solanki, and Unruh, 2000; Krivova and Solanki, 2008; Domingo et al., 2009; Shapiro et al., 2011a) and calculate the time-dependent solar spectrum as a sum of the spectra from the quiet Sun and different active features. We employ a four-component model
Figure 7 The implementation of the linear regression (red lines) through the four-orbit intervals of the LYRA data for the 13 January – 20 January 2010 period.

which treats separately contributions from the quiet-Sun, sunspots, active network, and plage areas. According to this model, the solar spectrum \( I(\lambda, t) \) can be written as

\[
I(\lambda, t) = \sum_k \left( \alpha_{QS}(\mu_k, t) I_{QS}(\lambda, \mu_k) + \alpha_S(\mu_k, t) I_S(\lambda, \mu_k) + \alpha_{AN}(\mu_k, t) I_{AN}(\lambda, \mu_k) + \alpha_P(\mu_k, t) I_P(\lambda, \mu_k) \right),
\]

(1)

where \( \alpha_{QS}(t) \), \( \alpha_S(t) \), \( \alpha_{AN}(t) \), and \( \alpha_P(t) \) are the time-dependent filling factors of the quite Sun, sunspots, active network, and plage areas, accordingly. \( I_{QS}(\lambda) \), \( I_S(\lambda) \), \( I_{AN}(\lambda) \), and \( I_P(\lambda) \) are the corresponding synthetic spectra. The summation represents the division of the solar disk in several concentric rings and is done over the different heliocentric angles, \( \mu_k \) is the corresponding cosines. Following the algorithm of Shapiro et al. (2011c), we used thirteen rings. This provides us with an accuracy of the order of a hundredth of a percent.
The synthetic spectra $I_{QS}(\lambda, \mu_k)$, $I_5(\lambda, \mu_k)$, $I_{AN}(\lambda, \mu_k)$, and $I_P(\lambda, \mu_k)$ are calculated with COSI. The temperature and density structure of the corresponding components are taken from Fontenla et al. (1999). A self-consistent simultaneous solution of the radiative transfer and the statistical equilibrium equation for the level populations guarantees that COSI considers the correct physics for the Herzberg region where the assumption of local thermodynamical equilibrium breaks down. The calculations with COSI yield the spectral solar irradiance which agrees well with the SOLSPEC measurements during the ATLAS 3 mission (Thuillier et al., 2011).

The time-dependent filling factors were extracted from the Solar Radiation Physical Modeling (SRPM: Fontenla et al., 1999, 2009) image mask of the Precision Solar Photometric Telescope (PSPT: Rast, Ortiz, and Meisner, 2008). In Figure 8 we give the dependency of the total (summed over all $\mu_k$ values) filling factors of the bright plage, plage, umbra, and penumbra on time. The profile of the dependencies represents the appearance and disappearance of the active features on the solar disk due to rotation (Figure 1) as well as the projection effect.

The comparisons of the LYRA data with SOLSTICE measurements and calculations are presented in the upper panel of Figure 9. The SOLSTICE data (available with 1-nm spectral resolution) and the calculated irradiance (available with 0.5 pm spectral resolution) were converted with the combined profile of the Herzberg filter and detector (Benmoussa et al., 2009).

Our four-component model of the solar variability is based on the models of the different solar-atmosphere components from Fontenla et al. (1999) who do not distinguish between the bright plage and plage as well as between the umbra and penumbra. At the same time...
the PSPT filling factors are based on the work of Fontenla et al. (2009) where these models are treated separately. To take this into account, we decreased the contrasts of the plage and sunspot with the quiet Sun in a way that the calculated variability in the Herzberg channel matches the SOLSTICE observations.

The theoretical understanding, as modeled by COSI, suggests that the variability in the Herzberg channel during the considered period has a one-peak profile. This can be explained by the dependency of the contrast between the different components of the solar atmosphere on the wavelength and heliocentric angle. The contrast between the bright components of the solar atmosphere (plage and active network) and quiet Sun strongly increases towards shorter wavelengths, while the contrast between the quiet Sun and sunspot depends on the wavelength more gradually (Figures 11 and 12 from Shapiro et al., 2010). As a result, for the considered transit of the active regions the increase of the Herzberg irradiance due to the presence of the plage and active network exceed the decrease of the Herzberg irradiance due to the presence of the sunspot independent of their position on the solar disk (for a more detailed explanation, see Unruh et al., 2008). The difference between LYRA and SOLSTICE or COSI indicates the uncertainty in the current processing of LYRA data.

4. Conclusions

The LYRA data in the Herzberg channel of the nominal unit are very much disturbed by the degradation and inhomogeneous flat field of the detector. However with careful and laborious analysis the real solar signal can be extracted. We believe that the processing of the data presented in this article is robust. We should also note that the main objective of this article should be seen not in the physical analysis of the solar variability (which is hindered by the aforementioned problems) but rather in the presenting the algorithm for the extracting the signatures of the solar variability from such noisy time series. Even if the LYRA data are not used in future, this algorithm can to some extent be useful for the analysis of data sets with similar problems.

The solar signal extracted from the LYRA data is in reasonable agreement with SORCE/SOLSTICE measurements (Figure 9). However the LYRA measurements indicate significant increase of the irradiance on 13 and 14 January, while there is no such increase in the SOLSTICE data. The theoretical results agree better with SOLSTICE measurements and yield a one-peak profile of the variability for the time period under consideration.

We believe that the origin of this disagreement can be clarified analyzing the data from the other recent European mission Picard launched on 15 June 2010. The PREMOS package onboard Picard consists of two experiments: one observing solar irradiance in five (two UV, one visible, and two near infrared) spectral channels with filter radiometers, the other measuring TSI with absolute radiometers. One of the PREMOS channels also measures the solar irradiance in the Herzberg continuum range (190 – 222 nm) but has a different type of detector (Schmutz et al., 2009). The intercomparison of the PROBA2/LYRA and Picard/PREMOS results will be addressed in a forthcoming article.

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