Photometric study of the star with changing Blazhko effect: TV Bootis

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ABSTRACT
This study investigates periodic modulation of the light curve of the RRc star TV Boo and its physical parameters based on photometric data. This phenomenon, known as the Blazhko effect, is quite rare among RRc stars. The frequency analysis based on the data gathered at the Masaryk University Observatory (MUO) and also using SuperWASP data revealed symmetrically structured peaks around the main pulsation frequency and its harmonics, which indicate two modulation components of the Blazhko period. The main modulation periodicity was found to be 9.7374 ± 0.0054 d. This is one of the shortest known Blazhko periods among RRc stars. The second modulation period (21.5 ± 0.2 d) causes changes of the Blazhko effect itself. Some indices show that TV Boo could be affected by long-term changes of the order of years. Basic physical parameters were estimated via an MUO V light-curve solution using Fourier parameters. TV Boo seems to be a low-metallicity star with [Fe/H] = −1.89.

Key words: stars: horizontal branch – stars: individual: TV Bootis – stars: variables: RR Lyrae.

1 INTRODUCTION
Light-curve modulation together with cyclic period changing is quite common among pulsation variables. In the case of RR Lyrae type stars, such behaviour has been known for more than one century and it is called the Blazhko effect after S. N. Blazhko, who first noticed phase modulation in RW Dra (Blazhko 1907). The amplitude modulation was discovered several years later by Shapley (1916) in the prototype RR Lyr itself. Nowadays, more than 40 per cent of all RRab stars seem to show this effect (Kolenberg et al. 2010; Jurcsik et al. 2006), but only about 5 per cent of RRC type (Moskalik & Poretti 2002).

Although many theories have been suggested to explain the observed properties of modulated stars, the origin of the Blazhko effect is still not well understood. A brief overview of these theories can be found in Smith (1995), Kovács (2009) and Kolenberg (2011). In the era of the Kepler space telescope, the 9:2 resonance between the fundamental radial mode and the ninth overtone seems to be responsible for the Blazhko effect (Buchler & Kolláth 2011).

The Blazhko effect manifests itself as sidelobe frequencies around the main pulsation frequency and its harmonics in the frequency spectra. Also the modulation frequency itself and its harmonics are natural manifestations of the modulation (e.g. Szeidl & Jurcsik 2009; Benkő, Szabó & Paparó 2011). Frequencies are usually equally spaced and form doublets, triplets or higher multiplets.

The variability of TV Boo (α = 14h16m36.58s, δ = +42°21′35.69″, J2000) was first reported in Gutnick & Prager (1926). GCVS (Samus et al. 2012) gives the period \( P = 0.3125609 \text{ d} \) and spectral type A7–F2. Metallicity ([Fe/H] = −2.44) listed in Feast et al. (2008) makes TV Boo one of the stars with the lowest metallicity among RR Lyrae type stars.

The Blazhko effect in TV Boo has been known since 1965, when Detre (1965) found a Blazhko period of 33.5 d. Almost 10 yr later, Firmanyuk (1974) reported the Blazhko effect with a period of 16.14 d. More recently, using data of the Northern Sky Variability Survey (NSVS; Wozniak et al. 2004), Wils et al. (2006) found the Blazhko period to be 10 d.

2 OBSERVATION AND DATA REDUCTION
TV Bootis was observed on 23 nights during three seasons from 2009 to 2011 (for the complete observation log, see Table 1) with the 60 cm Newtonian telescope of the Masaryk University Observatory (MUO) in Brno, Czech Republic. This telescope is equipped with an ST-8 CCD with a KAF-1600 chip (FOV 17.0 arcsec × 11.3 arcsec). Observations were carried out in BV\(_R\) filters. In each of these passbands, more than 2000 data points were collected. The data reduction and differential photometry were performed using the c-MUNIPACK package.\(^1\) TYC 3038-955-1 and TYC 3038-1064-1 were used as comparison and check stars, respectively. Although both stars are significantly redder than TV Boo itself (which could slightly affect the light curve and mainly colour indices), these stars were chosen because of the lack of other suitable stars in the vicinity of TV Boo. Standard deviations in the difference between the comparison and check stars were better than 0.01 mag in all passbands. All the points were transformed to the standard BV\(_R\_) magnitudes

\(^1\) http://c-munipack.sourceforge.net/
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3 FREQUENCY ANALYSIS

Frequency analysis was performed using the MUO data set as well as data from the SuperWASP (Butters et al. 2010) and NSVS (Wozniak et al. 2004) surveys. As the main tool for the frequency analysis Period 04 (Lenz & Breger 2004) was used. For basic pulsation period determination, PERSEA software [made by G. Maciejewski on the basis of the Schwarzenberg-Czerny (1996) method] and a Matlab script based on non-linear least-squares method were used. Within their margins of error, both methods gave the same results, so the maxima timings can be expressed as

$$ \text{HJD} T_{\text{max}} = 245\,4922.2334(2) + 0.312\,5615(7) E_{\text{puls}}. $$

The analysis based on the MUO data set gave almost the same results as were obtained with the SuperWASP data. Detected frequencies in all studied data sets are listed in Table 3. Semi-amplitudes, S/N ratios for individual frequencies and their standard deviations from the results obtained from Monte Carlo simulations are also given in this table. Results for NSVS data are not listed, because only three frequencies ($f_0$, $2f_0$ and $f_0 + f_0$) were detected in this data set.

The frequency analysis showed two modulation components located symmetrically around the main pulsation frequency and its harmonics $k f_0$ up to the third order in the case of SuperWASP data and only one component related to the second modulation frequency ($f_0 + f_m = 3.244\,555 \, \text{c d}^{-1}$) was revealed in the MUO data set. We detected eight harmonics of the main pulsation frequency in all studied data sets.

The positions of sidelobe peaks with higher amplitudes ($k f_0 \pm f_0$) correspond to $f_0 = 0.102\,697 \, \text{cycles/day} (\text{c d}^{-1})$, which results in a 9.7374 \, \text{d modulation period} (Blazhko period). Except for the sidelobe frequencies, other equally spaced peaks were detected in SuperWASP data around the main frequency with values $f_0 + f_m = 3.245\,98$ and $f_0 - f_m = 3.15276 \, \text{c d}^{-1}$. The difference between these frequencies and the main pulsation frequency is 0.04661 \, \text{c d}^{-1}, which indicates the second modulation period of length 21.5 \, \text{d}. This additional period is usually explained as a changing Blazhko effect, which is also known in other stars, e.g. XZ Cyg (Lagaly et al. 2004), RV UMa (Hurt et al. 2008) or LS Her (Wils et al. 2008).

A similar frequency structure with separate modulation peak series was also revealed in the case of CZ Lac (Sódor et al. 2011) and V445 Lyr (Guggenberger et al. 2011). All named stars (except LS Her) are of the RRc type with such a frequency structure.

According to the frequency analysis, the maxima timings for the Blazhko effect can be expressed as

$$ \text{HJD} T_{\text{max Blazhko}} = 245\,4950.80(10) + 9.7374(54) E_{\text{Blazhko}}. $$

Phased MUO V and SuperWasp light curves look very similar in shape and amplitude. Therefore, we tried to combine these data sets into one to evaluate the result. Phased light curves were fitted with the linear combination of sines to get a mean light curve and finally

\[ f \]
to obtain the shift in magnitudes between these two data sets. The frequency analysis of such data set revealed five frequency peaks related to the second modulation frequency, which were not detected in both data sets when studied separately. SuperWasp data cover a much broader wavelength range than the Johnson V passband. In addition, the shape of SuperWasp and MUO V light curve can be slightly different; therefore, such a combination is not rigorous.

One of the recent highlights of RR Lyrae research is the discovery of period-doubling (Szabó et al. 2010) and high-order resonances (Kolláth, Molnár & Szabó 2011) as possible causes of the Blazhko effect. No signs of such behaviour were unveiled in the frequency spectra of TV Boo. This could either be caused by insufficient quality of the data for this purpose or first overtone Blazhko pulsators do not show such behaviour.

Another puzzle seen in many RRc stars (and Cepheids) observed with high-precision photometry is the presence of an additional frequency with a strangely repeating frequency ratio of 1.58–1.63 from star to star (see e.g. Moskalik et al. (2012)). Again, no frequency peaks with the given ratio or even close to this ratio were identified in studied data sets down to the level of residuals. In addition, the detection of such a frequency was complicated by the fact that in the area defined by the ratio fell aliases of the main pulsation frequency and also aliases related to modulation frequencies. In any case, based on our analysis, we could not confirm or disprove the presence of the frequency in the range of 1.58–1.63. The fact that such a feature could not be confirmed might indicate that the feature is absent in the case of TV Boo.

Frequency spectra together with the spectral window resulting from the SuperWASP data set are shown in Fig. 2. The first two panels show Fourier spectra with the main pulsation frequency (a) and spectrum after removal of this frequency and its harmonics (b). There are easily noticeable daily aliases of the highest peaks in these two panels. The final residual spectrum, after pre-whitening with all basic and modulation frequencies, is given in panel (d). The height (about 3 mmag) of the residual spectrum is probably caused by small night-to-night shifts in the data as well as by unequally spaced observations. There is no significant peak in the residual plot.

### Table 3. Frequencies found in MUO and SuperWASP data sets.

| Id | Frequency (c^−1) | MUO A_V (mag) | S/N | Frequency (c^−1) | SuperWASP A (mag) | S/N |
|----|-----------------|-------------|----|-----------------|------------------|----|
| f_0 | 3.199 3740(8) | 0.2620(3) | 163.9 | 3.199 3646(4) | 0.2497(3) | 326.0 |
| f_0 + f_0 | 3.302 0366(6) | 0.0342(6) | 21.5 | 3.302 0663(3) | 0.0318(3) | 42.0 |
| f_0 − f_0 | 3.096 71(2) | 0.0109(4) | 6.8 | 3.096 67(1) | 0.0093(3) | 11.9 |
| f_0 + f_4 | 3.244 5553(3) | 0.0066(6) | 4.1 | 3.245 98(2) | 0.0060(3) | 7.9 |
| f_0 − f_4 | 3.152 76(2) | 0.0070(3) | 9.1 | 3.152 76(2) | 0.0069(3) | 11.2 |
| f_2 | 6.398 7484(4) | 0.0730(5) | 49.1 | 6.398 7382(2) | 0.0729(3) | 156.6 |
| f_2 + f_2 | 6.501 41(2) | 0.0133(5) | 9.0 | 6.501 43(2) | 0.0109(3) | 29.3 |
| f_2 − f_2 | 6.296 0693(2) | 0.0072(5) | 4.8 | 6.296 0693(2) | 0.0070(3) | 11.9 |
| f_4 | 9.700 78(3) | 0.0102(4) | 5.7 | 9.700 80(1) | 0.0077(3) | 18.2 |
| f_4 + f_4 | 9.645 34(3) | 0.0072(5) | 4.8 | 9.645 34(3) | 0.0070(3) | 11.9 |
| f_4 − f_4 | 9.598 12(1) | 0.0187(4) | 14.8 | 9.598 12(1) | 0.0185(3) | 43.6 |
| f_6 | 12.797 50(2) | 0.0088(4) | 9.1 | 12.797 50(2) | 0.0086(3) | 18.2 |
| f_6 + f_6 | 12.900 16(4) | 0.0051(4) | 5.2 | 12.900 20(3) | 0.0052(3) | 18.2 |
| f_6 − f_6 | 12.694 80(6) | 0.0021(3) | 5.7 | 12.694 80(6) | 0.0021(3) | 18.2 |
| f_8 | 15.996 87(2) | 0.0090(4) | 11.7 | 15.996 87(2) | 0.0089(3) | 22.5 |
| f_8 + f_8 | 16.043 41(7) | 0.0016(3) | 4.6 | 16.043 41(7) | 0.0016(3) | 4.6 |
| f_8 − f_8 | 15.849 62(5) | 0.0037(4) | 7.9 | 15.849 62(5) | 0.0036(3) | 11.2 |

4 Zero-points and amplitudes of the light changes were 11.123 and Δm 0.534 mag in the case of SuperWASP data and 10.917 and Δm 0.567 mag for MUO V data, respectively. The shift was therefore 0.206 mag.

### 4 Remarks on the Period Change During the Blazhko Cycle and Possible Long-Term Features

Blazhko RRc stars primarily show frequency modulation with only small amplitude modulation. In TV Boo, the amplitude of the light-curve modulation is only about 0.1 mag and the maximum and minimum of light changes are affected to a similar extent and the fuzziness of the light curve is uniform in all its parts (check Fig. 1).

TV Boo, as a frequency-modulated star with only small amplitude modulation, shows changes in the pulsation period during the Blazhko cycle. Changes in phase and amplitude of maximum light during three unique Blazhko cycles between HJD 2454151 and 2454181 can be seen in Fig. 3. Points in this plot represent maxima of our light-curve model. Points from the first and third Blazhko cycles create curves of almost the same shape, but the curve based on points coming from the cycle between them differs. The phase–amplitude diagram will be similar for a few following Blazhko cycles before it will change considerably. Bear in mind that the shape of the light curve is slowly changing with time, so the phase–amplitude diagram will look completely different after many Blazhko cycles.

Fig. 4 shows what the light curve looked like during one Blazhko cycle between HJD 245 4151 and 245 4161. Each Blazhko cycle
affects the light-curve shape uniquely in a given Blazhko phase, which is the manifestation of a changing Blazhko effect. Notice the change of light-curve shape around maxima during the Blazhko cycle.

The amplitude of the period changes determined from maxima timings is about 0.13 $P_{\text{puls}}$; in terms of time it is about 58 min. This value is strongly affected by the change of the light-curve shape around maxima. A more realistic estimation of the magnitude of the phase modulation can be obtained by measuring the phase shift in the ascending branch. We measured the phase shift in fixed magnitude 11.1 (approximately half of the light-curve amplitude) and we obtained the strength of period change during the Blazhko cycle beginning at HJD 245 4151 shown in Fig. 4 as 0.046 $P_{\text{puls}}$, which gives a little less than 21 min. This value is also slightly cycle-to-cycle dependent, but only of the order of 10$^{-3}$ $P_{\text{puls}}$.

Descending branches of the light curves are deformed with humps in some minima of the Blazhko cycles, but they are not present in each Blazhko cycle. In the MUO data, the hump was not detected in the season 2009, but it is clearly visible in the data gathered in 2010 and 2011. Similarly, in the SuperWASP data from the season 2007, this behaviour was not unveiled, but there is a weak sign of the hump in the 2004 season (see Fig. 5, middle panel). In the same figure, it is apparent that the Blazhko light curve is strongly asymmetric with its minimum around the Blazhko phase 0.7. This feature could be a demonstration of long-term changes known in few other examples, e.g. in RR Lyrae itself with long changes lasting 4 yr (Dreizler & Szeidl 1973). A first-order estimation of such a period could then be about 6 yr. However, this claim is based on the assumption that MUO $V$ and SuperWASP light curves are of the same shape and amplitudes. Of course, the absence of the hump can be caused by the properties of the data itself – lack of data for such Blazhko phases in studied data sets or the overlapping of data by other data.

5 PHYSICAL PARAMETERS

In the 1990s, some empirical relations between the so-called Fourier coefficients and basic physical parameters of RR Lyraes were found, namely there are calibrations for metallicity, temperature, absolute magnitude, etc. (Jurcsik 1998; Kovács 1998). The method is based on the decomposition of a phased $V$ band light curve into its sum of sines or cosines. This can be done for stable RR Lyraes and even for Blazhko stars, but good coverage of the folded light curve is needed. After decomposition, Fourier parameters $R_{ij}$ and $\phi_{ij}$ can be determined as follows:

$$V = A_0 + \sum_{i=1}^{n} A_i \sin(i\omega t + \phi_i),$$

$$\phi_{ij} = j\phi_i - i\phi_j,$$

$$R_{ij} = A_i/A_j.$$
Figure 4. Light-curve changes during one unique Blazhko cycle between HJD 245 4151 and 245 4161. The modelled light curve belonging to the given Blazhko phase (mentioned in the top-right corner of each plot) is plotted with a continuous line. The triangles show real SuperWASP data for the given Blazhko phase in this unique time span. All available SuperWASP data (plotted with dots) are just for illustration to see the range of possible light curves.

Figure 5. MUO V data (top panel), SuperWASP data from season 2004 (middle panel) and from season 2007 (bottom panel). Data folded with the Blazhko period are on the left-hand panels of the picture; data phased with the main pulsation period are to the right. Note the hump on the descending branch of MUO data and SuperWASP data (marked by an arrow).

Table 4. Fourier parameters based on sine-term decomposition.

| A_0 (mag) | R_{21} (mag) | R_{31} (mag) | R_{41} (mag) | \phi_{21} (rad) | \phi_{31} (rad) | \phi_{41} (rad) |
|-----------|--------------|--------------|--------------|----------------|----------------|----------------|
| 10.917    | 0.283        | 0.078        | 0.043        | 2.698          | 5.545          | 1.996          |

Table 5. Mean physical parameters of TV Boo.

| Parameter | Value       |
|-----------|-------------|
| [Fe/H]    | −1.89 ± 0.14|
| T_{eff} (K) | 7270 ± 50   |
| M_V (mag)  | 0.59 ± 0.06 |
| L (L☉)     | 56 ± 5      |
| R (R☉)     | 4.8 ± 0.5   |
| M (M☉)     | 0.73 ± 0.04 |
| r (pc)     | 1150 ± 30   |

used:

\[
[\text{Fe/H}] = 52.466 P^2 - 30.075 P + 0.131 (\phi_{31}^c)^2 + 0.9825 \phi_{31}^c - 4.198 \phi_{31}^c P + 2.424 \quad (6)
\]

\[
M_V = 1.061 - 0.961 P - 4.447 A_4 - 0.044 \phi_{21}^c \quad (7)
\]

\[
\log(L/L☉) = 2.41 + 1.04 \log P - 0.058 \phi_{31}^c \quad (8)
\]

\[
\log T_{eff} = 3.7746 - 0.1452 \log P + 0.0056 \phi_{21}^c \quad (9)
\]

\[
\log(M/M☉) = 0.39 + 0.52 \log P - 0.11 \phi_{31}^c \quad (10)
\]

where \( \phi_{ij}^c \) and \( \phi_{ij}^s \) are coefficients based on cosine-term decomposition and sine-term decomposition coefficients, respectively. The relations between sine and cosine coefficients are \( \phi_{ij}^c = \phi_{ij}^s - (j - k) \frac{\pi}{2} \). The zero-point in equation (7) was decreased according to Cacciari et al. (2005) by 0.2 as opposed to the equation in Kovács (1998).

Physical parameters resulting from equations (6)–(10) using the MUO V light-curve fit noted in Table 4 are listed in Table 5. Metallicity in equation (6) is given in the Zinn & West (1984) scale. The solution, and mainly parameter \( \phi_{31}^c \), is very sensitive to the coverage of the light curve. If we remove three nights when the hump was observed, the \( \phi_{31}^c \) decreases its value by about 0.04, which leads to a change in [Fe/H] of about −0.4. Our light curve is not ideally covered in all possible Blazhko phases, so we suspect metallicity to be even slightly higher (of the order of hundredths) than the obtained value [Fe/H] = −1.89.

This value differs by an order of a few tenths from the published values: [Fe/H] = −2.5 (Butler et al. 1982), [Fe/H] = −2.22 (Suntzeff et al. 1994), [Fe/H] = −2.44 (Fernley et al. 1998). These values came from the \( \Delta S \) approach (Preston 1959). In the case of the
Blazhko star TV Boo, we need to measure $\Delta S$ in various Blazhko phases to obtain a meaningful value of $\Delta S$. The problem of previous studies is that $\Delta S$ was not determined at minimum light, as it is required. Furthermore it was determined only in a few very close pulsation phases.

Published values of $\Delta S$ vary between 8 (Preston 1959) and 13.1 (Butler 1975). Transforming them via equation 3 in Suntzeff et al. (1994), we obtained a range of metallicities from $-1.67$ (Preston 1959) to $-2.44$ (Butler 1975). Solving these 10 light curves assuming one unique Blazhko cycle as shown in Fig. 4, we obtained metallicities between $-1.68$ and $-2.0$, so the determination of metallicity strongly depends on the Blazhko phase.

Peña et al. (2009) observed TV Bootis and determined basic physical parameters in the same way as we did. They proposed the metallicity to be $[\text{Fe}/\text{H}] = -2.04$. Unfortunately, their measurements cover only one pulsation cycle. The errors of physical parameters were calculated according to errors of calibrations given in papers they are taken from. All parameters (except metallicity) of TV Bootis in Table 5 are almost the same as found in the literature, especially with those listed in Peña et al. (2009). The value of the radius of star is only a rough estimation determined by solving the Stefan–Boltzmann equation for a blackbody.

The estimated mass of TV Boo is quite high compared to RRab type stars, but it is just in the middle of the mass interval for RRc stars noted in Simon & Clement (1993). The distance $r$ was determined using the distant modulus with extinction 0.024 mag in the V band. This value is in agreement with 1149 pc noted in Liu & Janes (1990).

6 CONCLUSIONS

The Blazhko star TV Bootis was analysed using data gathered at MUNI in 23 nights together with data from the SuperWASP survey. The data cover few years and contain a wide range of Blazhko phases. Therefore, we were able to determine new pulsation properties with the best precision ever and also we were able to obtain meaningful values of physical parameters.

The frequency analysis revealed symmetric structures located in the vicinity of the basic pulsation frequency and its harmonics with two modulation peaks. The star with the basic pulsation period a little bit shorter than 8 h shows light-curve modulation with a period of 9.7374 d. This is one of the shortest Blazhko period known for RR Lyrae type stars. The frequency analysis also unveiled another modulation period with a value of 21.5 d, which causes changes in the Blazhko effect itself. In addition to this period, there are some indications for longer light-curve changes of the order of years. This behaviour should be confirmed or disproved by further observations. No signs of the period-doubling phenomenon or a strange frequency in the range 1.58–1.63 $f_0$, revealed in some RRc stars, were detected.

We found TV Bootis to be a low-metallicity star with $[\text{Fe}/\text{H}] = -1.89$. This value is slightly higher than metallicities found in the literature (from $-2.5$ to $-2.04$). These values come from only few measurements based mainly on spectroscopic observations with a short time span, carried out in one unique Blazhko phase. The lack of observations in different Blazhko phases and low number of observations are a crucial deficiency in many studies of Blazhko stars. The discrepancies in metallicities could be resolved by a new spectroscopic measurements made in many Blazhko phases with sufficient time span.

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\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure6}
\caption{Phased MUO V light curve together with our solution. The data are plotted as points, and the fit with continuous lines.}
\end{figure}

\footnote{Taken from \url{http://ned.ipac.caltech.edu/forms/calculator.html}.}
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SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article:

Table 2. MUO V data from seasons 2009–2011 (http://mnras.oxfordjournals.org/lookup/suppl/doi:10.1093/mnras/sts135/-/DC1).

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