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High-Energy Photon and Particle Effects on Exoplanet Atmospheres and Habitability

Thematic Areas:
- Planetary Systems
- Star and Planet Formation
- Formation and Evolution of Compact Objects
- Cosmology and Fundamental Physics
- Stars and Stellar Evolution
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Abstract: It is now recognized that energetic stellar photon and particle radiation evaporates and erodes planetary atmospheres and controls upper atmospheric chemistry. Key exoplanet host stars will be too faint at X-ray wavelengths for accurate characterization using existing generation and future slated X-ray telescopes. Observation of stellar coronal mass ejections and winds are also beyond current instrumentation. In line with the Committee on an Exoplanet Science Strategy recognition that holistic observational approaches are needed, we point out here that a full understanding of exoplanet atmospheres, their evolution and determination of habitability requires a powerful high-resolution X-ray imaging and spectroscopic observatory. This is the only capability that can: (1) characterize by proxy the crucial, difficult to observe, EUV stellar flux, its history and its variations for planet hosting stars; (2) observe the stellar wind; (3) detect the subtle Doppler signatures of coronal mass ejections.
1 What Conditions Control Exoplanet Habitability?

The rate at which gas is lost from an exoplanet’s atmosphere is critical for the survivability of surface water. Atmospheric mass loss can be driven by both thermal and non-thermal processes, which depend upon the radiation and winds of their host stars. The dominant thermal process is hydrodynamical outflow energized by extreme ultraviolet (EUV; 100–912 Å) and X-radiation (0.1–100 Å) that heat the exoplanet’s thermosphere and levitate gas against the exoplanet’s gravitational potential (e.g. Owen & Jackson, 2012). Photodissociation and ionization of molecules, including water and CO$_2$, by the stellar UV and EUV radiation increases the mass-loss rate by producing lighter atoms (e.g., H) that are more easily lost to space. Most of the thermospheric heating is by EUV photons but this radiation cannot be observed directly because of interstellar H absorption. The chromospheric UV and FUV are inadequate EUV proxies. The strength and spectral energy distribution of a star’s EUV emission instead arises from the transition region and corona. The 30–60 Å range contains many of the same ionization stages that are important in the EUV range. Observing these enables prediction of the EUV spectrum. Detecting the relevant lines in exoplanet hosts requires a high-resolution ($R \geq 5,000$) spectrum that is not feasible with any existing or slated future missions, including Chandra, XMM-Newton or ATHENA.

The irradiation history of a planet also depends on the host star’s rotation rate, faster rotators producing larger radiation doses over time by an order of magnitude or more than slower rotators (Johnstone et al., 2015). To understand the range and likely radiation doses, it is essential to map out the EUV radiation through time for stars of similar ages but different rotation rates. This requires observations of open clusters with known ages at high spectral resolution in the soft X-ray range (30–100 Å) and a facility with with effective area of about 50× that of Chandra.

The X-ray emission of stars is variable on many time scales especially for M dwarfs, which many astronomers think are the best host star candidates for locating nearby habitable exoplanets. Young rapidly-rotating stars have high X-ray and EUV emission and emit energetic flares. Long-duration monitoring of the optical radiation of G-type stars by Kepler shows that high-energy superflares (total energy $E > 10^{32}$ ergs) are likely on a time scale of $\sim 500$ days for slowly rotating solar-like stars but are far more common on young G-type stars, and occur as often as 1 per 10 days (Shibayama et al. 2013). Superflares have been observed with energies as large as $E = 10^{35}$ ergs. Chandra has observed superflares on M dwarf and young stars, but the high-resolution spectra of superflares and also of more modest flares needed to infer their EUV emission require a high resolution large effective area soft X-ray spectrometer (see Figure[1]).

2 Stellar Winds and Exoplanet Atmospheric Loss

The flow of ionized stellar wind electrons and protons erode an exoplanet’s atmosphere. Ions produced by photoionization or charge-exchange reactions in the outer atmospheres of exoplanets can be picked up by the magnetic field in the stellar wind and expelled, can be lost through a
“polar wind”. Simulations show that such wind- and photoionization-driven processes can be a very important mass-loss agent for Earth-like planets around M stars (Garraffo et al., 2016; Dong et al., 2017; Garcia-Sage et al., 2017; Airapetian et al., 2017). Recent measurements by the MAVEN satellite (Brain et al., 2016) confirm previous estimates that the primary mass-loss mechanism for water on Mars is erosion by the solar wind.

The mass loss rates for late-type dwarfs are extremely difficult to measure as the solar mass-loss rate is only about $1.5 \times 10^{-14} M_\odot$ yr$^{-1}$. Radio observations yield only upper limits. There are indirect estimates of mass-loss rates up to 100 times larger for four G and K stars with stronger magnetic fluxes than the Sun based on Lyα absorption in the “wall” of hydrogen at the stellar analogy of the heliopause (Wood et al., 2014). There are only two estimates using this technique of mass-loss rates for M stars—$8 M_\odot$ for the active M3.5 dwarf EV Lac and an upper limit of $< 10 M_\odot$ for Proxima.

There is a clear need for new techniques for measuring the winds of a much larger sample of exoplanet host stars. Such a technique is enabled by sensitive, high spatial resolution X-ray imaging.

The ionized stellar wind interacts with neutral atoms in the ISM and the atmosphere through radiationless collisional transfer of one or sometimes multiple electrons from a neutral ISM atom or molecule to a wind ion. Electrons captured into the upper levels of highly ionized metals cascade to lower levels, emitting X-rays. The resulting X-ray spectrum is dominated by K-shell emission from H-like and He-like ions of C, O, N, and Ne. The conversion to wind mass loss rate is direct. An attempt by Wargelin & Drake (2002) to detect the charge exchange wind signature of Proxima using Chandra observations yielded only an upper limit of $3 \times 10^{-13} M_\odot$ yr$^{-1}$. Sub-arcsecond spatial resolution, high sensitivity and low background are required to make detections. With new detector technology such as the X-ray microcalorimeter, a sensitive next generation X-ray mission with arcsecond or better imaging will be able to observe the charge exchange signatures of stars out to at
least 10pc for solar-like mass loss rates, and to larger distance for higher rates, enabling winds to be mapped out with stellar activity level and spectral type and generally applied to exoplanet systems. Coronal plasma that is not confined by strong magnetic fields must participate in the stellar wind expansion. A mission with high soft X-ray resolution reaching $\lambda/\Delta\lambda = 5000$, corresponding to 60 km s$^{-1}$, and the possibility of measuring flow velocities three times smaller for bright emission lines, will also have the capability to measure stellar winds directly. This would be totally new science that only a large area, high resolution X-ray mission could accomplish.

3 Coronal Mass Ejections

Strong X-ray flares on the Sun are usually accompanied by the ejection of cooler material (roughly 10,000 K) that had previously been confined by magnetic fields that became disrupted during the flare. The ejected material, generally called coronal mass ejections (CMEs), may also contain high energy protons accelerated in the flare and CME shock front. CMEs differ from the quasi-steady solar wind in two respects: they are orders of magnitude denser, and are spatially confined.

Segura et al. (2010) modeled the effect of a superflare ($E \approx 10^{34}$ erg) and CME impact on a hypothetical Earth-like exoplanet located in the habitable zone (0.16 AU) of the flare star AD Leo (dM3e). High energy protons with energies greater than 10 MeV severely depleted nitrogen oxides, and subsequently ozone, in the atmosphere for 2 years. Airapetian et al. (2016) found CME energetic particles can create important prebiotic molecules and alter atmospheric greenhouse gases potentially important for the Faint Young Sun paradox.

These studies demonstrate the acute need for observations of stellar CME events. No such events have been definitively detected, although there are searches underway at low frequency radio wavelengths. Extrapolations of solar CME-flare relationships (Figure 2) are uncertain by orders of magnitude but are sorely needed to understand what CME activity exoplanets are experiencing. High-energy protons are very difficult to observe, but the cooler material in stellar CMEs, or the associated compression wave in the corona, should be observable by a sensitive high resolution X-ray spectrometer. There are two X-ray detections of probable CMEs where the cool, dense material is seen in absorption as it passes in front of the flaring corona: The 20 August 1980 flare on Proxima Cen observed by Einstein Haisch et al. (1983); and the 30 August 1997 superflare on Algol observed by BeppoSAX (Moschou et al., 2017).

High-resolution spectroscopy at X-ray wavelengths could routinely and definitively observe the tell-tale Doppler shifts of CMEs or their coronal compression waves (Figure 2) and identify their physical properties, including their thermal structure, masses and energies. A combination of high throughput and high spectral resolution will be critical, mapping out CME frequency and energy vs optical and X-ray flare diagnostics for exoplanet hosts directly, and generally as a function of spectral type and activity level.

4 Transmission spectroscopy of exoplanet atmospheres

X-rays are powerful diagnostics of planetary upper atmospheric gas density structure and chemical composition. The transit of the hot Jupiter HD189733b was detected through X-ray absorption by oxygen in Chandra observations by Poppenhaeger et al. (2013), who found the scale height of
X-ray absorbing gas was higher than suggested by optical and UV transits. Hot Jupiters and similar giant close-in planets are important for improving theory and models describing atmospheric loss. X-ray absorption measures gas bulk chemical composition (Figure 3) along the line-of-sight—in this case in the transiting exoplanet atmosphere backlit by the host star’s corona. Such measurements are unique to the X-ray range, but only the very closest hot Jupiters are accessible with Chandra and XMM-Newton, and then only at low signal-to-noise ratio. An observatory with an effective area $50 \times$ that of Chandra will be able to observe HD 189733b-like transits out to 140 pc, a factor of more than 300 improvement in survey volume over current missions. Combination with optical/IR data will provide a powerful probe for clouds and hazes that can confuse IR spectroscopic analyses (Sing et al., 2016). By coadding observations of many transits, such a mission could also open studies to larger habitable planets, such as super Earths around nearby M dwarfs (Figure 3).

5 Summary

Exoplanet atmospheric loss and evolution cannot be properly understood without a powerful X-ray observatory capable of high spectral resolution of $R \geq 5,000$ at soft X-ray wavelengths, a microcalorimeter for higher energy high resolution imaging spectroscopy, a large effective area at least several decades greater than that of Chandra, and with spatial resolution better than 1 arcsecond.
Figure 3: A large area X-ray observatory will be able to measure gas bulk composition from transmission spectroscopy. Left: Illustration of the enormous difference in X-ray transmittance of gas with solar and Earth’s atmosphere compositions. Right: Simulation of detection of the 0.5 keV oxygen absorption edge betraying enhanced O abundance for 100 transits of a superearth planet around an M dwarf for a telescope with $50 \times$ the area of Chandra (by K. Poppenhaeger).

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