On Online Energy Harvesting in Multiple Access Communication Systems

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Abstract—We investigate performance limits of a multiple access communication system with energy harvesting nodes where the utility function is taken to be the long-term average sum-throughput. We assume a causal structure for energy arrivals and study the problem in the continuous time regime. For this setting, we first characterize a storage model that captures the dynamics of a battery with energy harvesting and variable transmission power. Using this model, we next establish an upper bound on the throughput problem as a function of battery capacity. We also formulate a non-linear optimization problem to determine optimal achievable power policies for transmitters. Applying a calculus of variation technique, we then derive Euler-Lagrange equations as necessary conditions for optimum power policies in terms of a system of coupled partial integro-differential equations (PIDEs). Based on a Gauss-Seidel algorithm, we devise an iterative algorithm to solve these equations. We also propose a fixed-point algorithm for the symmetric multiple access setting in which the statistical descriptions of energy harvesters are identical. Along with the analysis and to support our iterative algorithms, comprehensive numerical results are also obtained.

Index Terms—Energy harvesting, Multiple access communication, iterative algorithm.

I. INTRODUCTION

The direct impact of energy on communication cost and lifetime has spurred significant efforts to manage and optimize energy consumption. In this respect, current and future state of the art technology has focused on harvesting energy from the environment. It is thus of paramount importance to design suitable adaptive power transmission policies for these technologies. In particular, the formulation of power policies in energy harvesting systems depends intricately on many factors, including energy arrival distribution, battery capacity, quality of service, etc. Moreover, most renewable energy resources have unpredictable behaviour that makes the design process of optimal power policies difficult. Solar panels, for instance, can only scavenge sunlight during the daytime and even then, this is a function of weather conditions. Another example is thermoelectric generators where energy is absorbed based on random temperature gradients between two metal junctions. Regarding these examples, a key objective of recent studies is to engineer optimal transmission power policies. These studies, depending on causal or non-causal knowledge of future energy arrivals, fall within two major categories: offline or deterministic (for non-causal), and online or stochastic (for causal) energy harvesting systems.

In the offline regime and in terms of throughput maximization, optimal power allocation for different communication topologies has been well studied. For instance, [2] studies the multiple access channel (MAC), [3] studies the broadcast channel, and the interference channel is studied in [4]. In addition, the issue of maximizing throughput in a fading channel has been treated in [5]. There, a directional water-filling algorithm is proposed. In [6], a continuous time energy harvesting system with constant energy leakage rate due to battery imperfections is considered. Another interesting problem has been studied in [7] where an offline energy harvesting problem subject to minimizing the transmission completion time is analyzed. Specifically, a continuous-time policy to minimize the delivery time of data packets is formulated. Among more recent results in the offline setting is [8] where energy cooperation in a two-hop communication system is considered.

As an overview of prior works in the online regime, we refer the reader to [5], [9], [10], and [11]. In [5] an algorithm in the offline problem of throughput maximization by a deadline was heuristically applied to the online counterpart. The authors have also considered a dynamic programming solution for online policies. However, the curse of dimensionality in the backward induction renders the computational cost of this approach very expensive. In [9], the capacity of the additive white Gaussian noise channel (AWGN) under discrete-time energy arrivals and infinite battery capacity is characterized. Additionally, two achievable schemes based on save-and-transmit and best-effort-transmit are studied there. In [10], queuing aspects of the online energy harvesting problem with infinite battery and buffer capacity have been considered. The authors have also suggested a greedy policy that in the low signal to noise ratio (SNR) regime is throughput optimal and attains minimum delay. A more relevant study related to the work presented here is [12]. Therein, Srivastava and Koksal have investigated an optimization problem where the objective is to maximize a utility function subject to causality and battery constraints. More interestingly, they addressed a trade-off between achieving the optimum utility and keeping the discharge rate low.

In this paper, we consider the online setting with continuous time policies in which the energy release rates are regulated dynamically based on the remaining charge of the battery at each moment. This architecture naturally requires a different mathematical framework in terms of modelling and analysis. Particularly, the main tool here for modelling the interaction between battery, energy arrivals, and energy consumption is a stochastic process known as a compound Poisson dam model. This model was pioneered by Moran in 1954 [13] and studied further by Gaver-Miller [14] and Harrison-Resnick [15]. In connection with this model, we derive an upper bound on the total sum-throughput of an online energy harvesting system. Also in terms of achievability, we construct an optimization problem to maximize the sum-throughput subject to an ergodicity constraint. This maximization problem turns out
to be non-linear and analytically cumbersome. Relying on a calculus of variations approach, we subsequently find a system of simultaneous PIDEs as necessary conditions for an optimal power policy. We then propose a Gauss-Seidel method (see [16]) to solve these equations efficiently. In the symmetric case, when the statistical description of all the energy harvesters are identical, we obtain an alternative algorithm using a fixed point iteration method. Moreover, in the case of the point-to-point channel setting, the necessary condition further reduces to a non-linear, autonomous ordinary differential equation (ODE) that can be solved directly, using conventional numerical methods [1].

The rest of the paper is organized as follows. In section II, we review some background, definitions, and notation. Section III deals with necessary and sufficient conditions for ergodicity of the storage process. In Section IV, we derive an upper bound as well as the achievable results for both finite and infinite storage cases, including two algorithms for the achievability part. These algorithms are then used to compute the numerical results in Section V. Lastly, in Section VI, we summarize our main findings and outline possible future directions.

II. PRELIMINARIES

A. Communication model

We consider $M$ multiple access transmission nodes that wish to transmit their data over a shared communication channel. Furthermore, each transmission node has an energy harvesting module and a battery to capture and store arriving energy packets. Throughout the paper, we denote the instantaneous transmission power at time $t$ from the $k$th node ($k = 1, 2, \ldots, M$) by $P_k(t)$. Also, to quantify the corresponding transmission rate of the nodes, we consider Shannon’s rate function, $r(x) = \frac{1}{2} \log_2 (1 + (x/N_0))$, where $N_0$ denotes the noise power spectral density. In particular, Shannon’s rate function carries the following properties and, unless stated otherwise, only these properties will be used in Section IV:

(A.1) Positivity: $r(x) > 0$ for all $x > 0$ and $r(0) = 0$.

(A.2) Differentiability: $r(x)$ is three times continuously differentiable on $x \geq 0$.

(A.3) Monotone increasing: $r'(x) > 0$ for all $x \geq 0$.

(A.4) Concavity: $r''(x) < 0$ for all $x \geq 0$.

Letting $R_k$ denote the long-term average rate of the $k$th user, we then have the rate-region described by

$$\sum_{k \in S} R_k \leq \lim_{T \to \infty} \frac{1}{T} \int_0^T r \left( \sum_{k \in S} P_k(s) \right) \, ds,$$

where the inequality holds for all subsets $S \subseteq \{1, 2, \ldots, M\}$, and the resulting region is a polytope called polymatroid. In this study, we restrict ourselves to the dominant face of this polymatroid (called permutahedron) that represents the total sum-throughput (or sum-rate) of the channel. Then, the sum-throughput is

$$\sum_{k=1}^M R_k = \lim_{T \to \infty} \frac{1}{T} \int_0^T r \left( \sum_{k=1}^M P_k(s) \right) \, ds.$$

B. Energy harvesting and storage model

In our energy harvesting model, we allow the transmission nodes to use different techniques for harvesting exogenous energy. For example, while one node may collect solar energy, another node can use a thermoelectric generator. This mechanism is especially important for sensor networks where distributed terminals may measure miscellaneous targets that also feed sensors with energy (e.g. see [17]). Mathematically, we assume that for each individual node $k \in \{1, 2, \ldots, M\}$, energy is rephased into the corresponding battery according to specific energy arrivals $E_k^n, E_k^{n+1}, \ldots$, where the superscript denotes the order of arrivals. Furthermore, the energy arrivals for node $k$ are independent, identically distributed (i.i.d.) according to $\mathbb{P} \{ E_k \leq x \} = B_k(x)$ which occur at random arrival times denoted by $T_k^n, T_k^{n+1}, \ldots$. The interarrival times $\Delta T_k^n = T_k^{n+1} - T_k^n$ are also assumed to be i.i.d. and exponentially distributed. Therefore, the attributed point process, $N_k(t) \overset{\text{def}}{=} \sum_{n \in \mathbb{N}} 1\{ \tau_k < t \}$, is a homogeneous Poisson point process with intensity denoted by $\lambda_k$. Consequently, the total energy flow $E_k^\text{in}(0, t)$ into node $k$ and up to time $t$ is a compound Poisson process,

$$E_k^\text{in}(0, t) \overset{\text{def}}{=} \sum_{i=0}^{N(t)} E_k^i.$$

To characterize the storage model, we also need to determine the output process at each transmitter. To do so, let $X_k(t)$ denote the energy stored in the $k$th battery as a function of time. Then, the total energy expenditure until time $t$ is

$$E_k^\text{out}(t) \overset{\text{def}}{=} \int_0^t P_k(s) \, ds,$$

where $P_k(\cdot)$ represents the transmission power policy of the $k$th transmitter, modulated by the available energy in the battery. Now, the storage equation in terms of the energy arrivals in Eq. (3) and the drift process in Eq. (4b) is

$$X_k(t) = X_k(0) + E_k^\text{in}(0, t) - \int_0^t p_k(X_k(s)) \, ds,$$

where $X_k(0)$ is the initial battery reserve at time $t = 0$, and here the battery is assumed to have infinite capacity ($X_k(t) \in [0, \infty)$). In the case that the $k$th battery has a finite storage capacity, say $L_k$, then $X_k \in [0, L_k]$, and we can similarly characterize the following dynamics,

$$X_k(t) = X_k(0) + E_k^\text{in}(0, t) - \int_0^t p_k(X_k(s)) \, ds - Z_k(t),$$

where $Z_k(t)$ is $\mathbb{R}^+$ valued process that is null at zero ($Z_k(0) = 0$), non-decreasing, continuous almost everywhere, and such that $\int_{\mathbb{R}^+} (L_k - X_k(s)) \, dZ_k(s) = 0$. This process, known as reflection process [18], ensures that for any energy arrival, the storage process remains inside the boundary, i.e., $X_k(t) \in [0, L_k]$.

It is also interesting to note that the application of the structures in Eqs. (5) and (6) are not limited to the current problem. In fact, this formulation has wide applicability in other fields of studies. Examples include workload modulated queues [19], water reservoir dam analysis [20], food contaminants exposure in bioscience [21], etc. In this paper, the ergodicity results of [20] will be used and are summarized in section III.

Notation. In the rest of the paper and for conciseness, we adopt several shorthand notations. In particular, $[M]$ stands for $\{1, 2, \ldots, M\}$. For $M > 1$, we define the rectangular domain $A$ as

$$A \overset{\text{def}}{=} [0, L_1] \times [0, L_2] \times \cdots \times [0, L_M].$$
Related to this, we also define the $M$ dimensional integral by
\[
\int_{L_1}^{L_2} \cdots \int_{L_M} \cdot d{x_1} d{x_2} \cdots d{x_M},
\]
which is represented by $\int_{A(S)} d{x}$. For all subsets $S \subseteq [M]$, we use $A(S)$ to denote the projection of $A$ onto the coordinates indexed by $S$, i.e.,
\[
A\{1, 3\} = [0, L_1] \times [0, L_3].
\]
Then, $\int_{A(S)} d{x}$ denotes integration over a subset of $\mathbb{R}^{|S|}$. $A_j$ is also a shorthand for
\[
A_j \equiv [0, L_1] \times \cdots [0, L_{j-1}] \times [0, L_{j+1}] \cdots [0, L_M].
\]

### III. Ergodic Theory of Storage Process

We here summarize necessary and sufficient conditions for ergodicity of the storage process in Eq. (5). Before stating the definitions regarding ergodic behaviour, we first put some mild constraints on the transmission policies. Particularly, for all $k = 1, 2, \ldots, M$,
1. $\forall L_k > 0$, $0 < x_k \leq L_k \Rightarrow p_k(x_k) > 0$ and $p_k(0) = 0$,
2. $\forall L_k > 0$, $\sup_{0 < x_k \leq L_k} p_k(x_k) < \infty$.

The first condition indicates that as long as there is energy in the battery, transmission continues (otherwise, the battery would have a minimum energy reserve that can not be consumed). The second condition does not permit the energy in the battery to be consumed instantly. Regarding these constraints, we say a policy is admissible if it fulfills these two conditions.

**Definition 1.** (Hitting Time) The hitting time, $\tau(x)$, is defined as the first time that the energy level in the battery reaches the value of $x$. More specifically,
\[
\tau(x) \equiv \inf\{t \geq 0 : X(t) = x\}.
\]

**Definition 2.** (Transient and Recurrent Process) The storage process is said to be transient, if and only if for all initial energy levels $x(0)$ in the battery, we have $\mathbb{P}\{X_t \to \infty\} = 1$. Alternatively, the storage process is said to be recurrent if and only if $\mathbb{P}\{\tau(x) < \infty|x(0)\} = 1$, $\forall x > 0$, $x(0) > 0$. In the case of a recurrent storage process, it is said to be positive recurrent if it further satisfies $\mathbb{E}[\tau(x)|x(0) = x] < \infty$ for one $x > 0$ and therefore for all $x > 0$ (irreducibility). Similarly, the recurrent storage process is null recurrent if $\mathbb{E}[\tau(x)|x(0) = x] = \infty$ for one $x > 0$ and therefore for all $x > 0$.

One motivation for surveying ergodic conditions is to rule out policies that result in transient and null recurrent battery behaviours. For example in the transient case $X(t) \to \infty$ a.s. which is unrealistic. Also, in the null recurrent case $\lim_{t \to \infty} \mathbb{P}\{X(t) \leq u|x(0) = x\} = 0$, $\forall x, u \geq 0$ which implies an unbounded energy reserve in the battery.

**Theorem 1.** (Ergodicity Condition) The storage process $\{X_k(t)\}_{t \geq 0}$ is positive recurrent if and only if there exist a probability measure $\pi_k$ that is absolutely continuous on $(0, \infty)$ and which may possess an atom at zero, $\pi_k^0 = \pi_k\{0\}$, i.e.,
\[
\pi_k(x_k) = \pi_k^0 + \int_{0^+}^{x_k} f_k(v_k) \, dv_k,
\]
and such that
\[
f_k(x_k) = \frac{\lambda_k}{p_k(x_k)} \left(p_k^0(1 - B_k(x_k)) + \int_{0^+}^{x_k} (1 - B_k(x_k - v_k)) f_k(v_k) \, dv_k\right). \quad (8)
\]
Furthermore, $\pi_k^0$ is the unique stationary distribution of the process $X_k(t)$.

**Remark 1.** The elegant proof of Assmussen for the converse part of Theorem 1 is based on an embedded Markov chain $\{X_k(n)\}$ at marked arrival times. In particular, for recurrent embedded chains, it is shown that any storage interval $[x_k^1, x_k^2)$, $0 < x_k^1 < x_k^2$ is recurrent in the sense of Harris. An alternative proof of the converse part of Theorem 1 adopts the additional condition $\int_0^\infty (1/p_k(u)) \, du < \infty, \forall x_k > 0$. Due to this extra condition, the required time to reach the zero state in the absence of new arrivals from any energy level in the battery must be finite. For this constraint, it can also be shown that $x_k = 0$ is a regenerative recurrent point for the process and therefore, due to the additional constraint, the probability measure has a strict atom $\pi_k^0 > 0$ at zero.

**Remark 2.** As discussed in [20, p. 297], in the finite energy case ($L_k < \infty$), the storage process is always positive recurrent and the probability measure is likewise governed by Eqs. (7) and (8).

**Remark 3.** We note that the atom of the probability measure $\pi_k(x_k)$ corresponds to an absorbing state of the process $X_k(t)$ in the sense that upon $X_k(t)$ entering state $x_k = 0$, the process remains there until an energy arrival occurs (at which point the process transits to another state). Based on this and the first constraint on admissible power policies (in particular $p_k(L_k) > 0$), there is no atom at $x_k = L_k$ in the finite case since it has a strictly negative drift in Eq. (6) that shifts the process to the inner region of the state-space instantaneously, i.e., $x_k < L_k$. Therefore, the battery never idles with $x_k = L_k$ (reflecting boundary).

An interpretation for the forward equation in Eq. (8) can be provided in terms of level crossing theory. In particular,
\[
f_k(x_k) p_k(x_k) = \frac{\lambda_k}{p_k(x_k)} \left(p_k^0(1 - B_k(x_k)) + \int_{0^+}^{x_k} (1 - B_k(x_k - v_k)) f_k(v_k) \, dv_k\right), \quad (9)
\]
is the equilibrium condition between the rate of down crossing at level $x_k$ (the l.h.s of Eq. (9)) and up crossing at level $x_k$ (the r.h.s of Eq. (9)). We can also view (8) as a Volterra integral equation of the second kind with the kernel function $K(x_k, v_k) = 1 - B_k(x_k - v_k)$, and it can thus be solved numerically (see [22]).

In this paper, we consider the energy arrivals $\{E_k\}_{k=0}^{\infty}, k = 1, 2, \ldots, M$, to be exponentially distributed with parameter $\zeta_k$. Therefore, we have
\[
K(x_k, v_k) = \exp(-\zeta_k(x_k - v_k)),
\]
that simplifies (8) to
\[
f_k(x_k) = \frac{\lambda_k \exp(-\zeta_k x_k)}{p_k(x_k)} G_k(x_k), \quad (10)
\]
where
\[
G_k(x_k) \equiv \left(p_k^0 + \int_{0^+}^{x_k} \exp(\zeta_k v_k) f_k(v_k) \, dv_k\right). \quad (11)
\]
Remark 4. The storage models in Eqs. (5) and (6) are memoryless, in the sense that at each available instant \( t \), the power policy \( p_k \) only depends on the available charge \( X_k(t) \) in the battery and not the entire sample path \( \{ X_k(s); s \leq t \} \). As an extension, we can also define a storage model with memory and infinite battery capacity as follows

\[
X_k(t) = X_k(0) + E_k^n(0,t) - \int_0^t p_k(X(u); u \leq s) \, ds. \tag{12}
\]

The extension of the storage model with memory and finite battery capacity follows similarly. However, when the arrival process is Poisson, it can be shown that \( X_k(t) \) is a sufficient statistic for an optimal power policy for both infinite and finite battery cases (see Appendix A). In this regard, knowledge of the entire path \( \{ X_k(s); s \leq t \} \) as an argument of \( p_k(\cdot) \) is excessive.

IV. BOUNDS ON TOTAL AVERAGE THROUGHPUT

Our objective now is to derive an upper bound on the average throughput as well as achievable policies with good performance. In connection with our system model, we will analyze a MAC with 1) finite, and 2) infinite storage batteries.

In particular, in the finite storage case, a good power policy must manage overflow in the battery as regular overflow causes energy waste and potentially decreases the sum throughput. To reduce overflow, the power policy must result in a large transmission power when the battery charge is large as otherwise overflow is likely to occur upon a new arrival. However, transmitting with too large a transmission power when the battery happens to have large charge is also undesirable due to the concavity of the rate function. In other words, there is a tension between overflow and the rate at which the large battery charge is consumed to reduce overflow likelihood.

To further clarify the latter point, consider an energy harvesting system with a single node \(( M = 1 \) in which energy \( E \) is replenished into a battery exactly every \( \alpha T \) units of time. In addition, assume that the transmitter sends data by using a constant transmission power \( P = E/(\alpha T) \), \( \alpha > 0 \). Two cases can now be examined:

(i) \( \alpha > 1 \): In this case, the transmitter fails to consume the entire battery charge before the next arrival, and thus overflow occurs regularly. We then have

\[
T \times r \left( \frac{E}{\alpha T} \right) \leq T \times r \left( \frac{E}{T} \right). \tag{13}
\]

(ii) \( \alpha < 1 \): In this case, the transmitter depletes its available battery charge within \( \alpha T < T \) of each arrival. From the concavity of the rate function, we have the following inequality

\[
\alpha T \times r \left( \frac{E}{\alpha T} \right) \leq T \times r \left( \frac{E}{T} \right). \tag{14}
\]

Here, the tension between (i) and (ii) is resolved by the optimal choice of \( \alpha = 1 \), i.e., \( P = E/T \).

A. An Upper Bound

1) Finite Storage Battery: In this case \( L_k < \infty \), \( \forall k \in [M] \). Then from (2) and due to ergodicity of the storage processes \( \{ X_k(t); t \geq 0 \} \) in the finite battery case (ref. Remark 3), we have almost surely

\[
\sum_{k=1}^M R_k \overset{a.s.}{=} \mathbb{E} \left[ r \left( \sum_{k=1}^M p_k(X_k) \right) \right], \tag{15}
\]

where the expectation is with respect to the stationary distribution in Theorem 1. In addition, from the concavity property of the rate function and Jensen’s inequality,

\[
\mathbb{E} \left[ r \left( \sum_{k=1}^M p_k(X_k) \right) \right] \leq r \left( \sum_{k=1}^M \mathbb{E}[p_k(X_k)] \right). \tag{16}
\]

It thus remains to bound the mean transmission power \( \mathbb{E}[p_k(x_k)] \). This can be accomplished by integrating by parts as follows

\[
\mathbb{E}[p_k(X_k)] = \pi_k^0 p_k(0) + \int_0^{L_k} p_k(x_k) f_k(x_k) \, dx_k \tag{17}
\]

\[
\overset{(a)}{=} \int_0^{L_k} p_k(x_k) f_k(x_k) \, dx_k \tag{18}
\]

\[
\overset{(b)}{=} \lambda_k \int_0^{L_k} \exp(-\xi_k x_k) G_k(x_k) \, dx_k \tag{19}
\]

\[
= -\frac{\lambda_k}{\zeta_k} \left[ \exp(-\xi_k x_k) G_k(x_k) \right]_0^{L_k} \tag{20}
\]

\[
+ \frac{\lambda_k}{\zeta_k} \int_0^{L_k} \exp(-\xi_k x_k) G_k'(x_k) \, dx_k,
\]

where (a) comes from the first constraint on the admissible power policies and (b) follows from (10). Now from (11),

\[
G_k'(x_k) = f_k(x_k) \exp(\xi_k x_k). \tag{21}
\]

Also we note that \( G_k(0^+) = \pi_k^0 \) and

\[
e^{-\xi_k L_k} G_k(L_k) = e^{-\xi_k L_k} \left( \pi_k^0 + \int_0^{L_k} \exp(\xi_k x_k) f_k(x_k) \, dx_k \right) \tag{22}
\]

\[
\geq e^{-\xi_k L_k} \left( \pi_k^0 + \int_0^{L_k} f_k(x_k) \, dx_k \right) \tag{23}
\]

\[
= e^{-\xi_k L_k}, \tag{24}
\]

where inequality (c) is due to the fact that \( \exp(\xi_k x_k) \geq 1 \) for all \( x_k \geq 0 \) since \( \xi_k > 0 \). Substituting (24) and (21) in Eq. (20) yields

\[
\mathbb{E}[p_k(x_k)] = \frac{\lambda_k}{\zeta_k} \left( G_k(0^+) - e^{-\xi_k L_k} G_k(L_k) \right) \tag{25}
\]

\[
+ \frac{\lambda_k}{\zeta_k} \int_0^{L_k} f_k(x_k) \, dx_k,
\]

\[
\leq \frac{\lambda_k}{\zeta_k} \left( \pi_k^0 - e^{-\xi_k L_k} + \int_0^{L_k} f_k(x_k) \, dx_k \right) \tag{26}
\]

\[
= \frac{\lambda_k}{\zeta_k} (1 - \exp(-\xi_k L_k)), \tag{27}
\]

In the last step, we now use (27) and the non-decreasing property of the rate function to characterize an upper bound for all \( L_k < \infty \) as follows

\[
\sum_{k=1}^M R_k \leq r \left( \sum_{k=1}^M \frac{\lambda_k}{\zeta_k} (1 - e^{-\xi_k L_k}) \right) \overset{\text{def}}{=} R_{\text{upper}}. \tag{28}
\]

2) Infinite Storage Battery: We now take \( L_k = \infty \). In this case, similar to (19) we can directly compute,

\[
\mathbb{E}[p_k(x_k)] = \lambda_k \int_0^{\infty} e^{-\xi_k x_k} G_k(x_k) \, dx_k \tag{29}
\]

\[
= \lambda_k \int_0^{\infty} e^{-\xi_k x_k} \left( \pi_k^0 + \int_0^{x_k} e^{-\xi_k v_k} f_k(v_k) \, dv_k \right) \, dx_k \tag{30}
\]
where in (a), we changed the order of integration. Thus, for positive recurrent policies and when all \( L_k = \infty \), we have the following upper bound

\[
\sum_{k=1}^{M} R_k \leq r \left( \sum_{k=1}^{M} \frac{\lambda_k}{\zeta_k} \right) .
\]  

(35)

Remark 5. In contrast with the inequality (28) which only holds for positive recurrent transmission power policies, Eq. (35) is valid for transient and null recurrent power policies as well. In particular, in the infinite battery case,

\[
\lim_{T \to \infty} \frac{1}{T} \int_{0}^{T} p_k(X_k(t)) \, dt \leq \lambda_k / \zeta_k ,
\]

regardless of the type of power policy, and thus (35) follows by concavity of the rate function. Nevertheless, the strict equality in Eq. (34) will be used to study transmission power policies that result in ergodic behavior for the infinite battery capacity case in Section IV-B.

B. Achievable allocation scheme

To derive transmission power policies with good performance, we start with the ergodicity assumption and the definition of expectation, i.e.,

\[
\sum_{k=1}^{M} R_k = \lim_{T \to \infty} \frac{1}{T} \int_{0}^{T} r \left( \sum_{k=1}^{M} P_k(s) \right) \, ds \quad \text{s.t.:} \quad \sum_{k=1}^{M} \frac{\lambda_k}{\zeta_k} \sum_{k=1}^{M} \pi_k(d x_k) = \bar{R} \left( \{ p_k(x_k) \}_{k=1}^{M} \right) ,
\]

(36)

(37)

(38)

\[ \pi_k(d x_k) = [\pi_k^0(\delta(x_k) + f_k(x_k))] \, d x_k , \]

(39)

and \( \delta(x_k) \) denotes the Dirac delta function. We now aim to find achievable policies through the following optimization problem

\[
\sup_{\{ p_k^0(x_k) \}_{k=1}^{M}} \int_{0}^{T} r \left( \sum_{k=1}^{M} p_k(x_k) \right) \prod_{k=1}^{M} \pi_k(d x_k) ,
\]

(40a)

(40b)

(40c)

(40d)

which maximizes the overall expected throughput of the multiple access channel subject to the stationary probability measure constraints of the batteries. However, tackling this non-linear optimization problem is challenging as the feasibility constraint in Eq. (40b) is not in an explicit form. To circumvent this difficulty, we use a calculus of variations approach to transform the problem into a set of necessary conditions for an optimal solution. As a starting point, consider the following linear mappings

\[
g_k(x_k) = f_k(x_k) e^{\zeta_k x_k} , \quad x_k > 0 ,
\]

(41)

that transforms the positive recurrent condition in Eq. (10) into

\[
g_k(x_k) = \frac{\lambda_k}{p_k(x_k)} \left( r_k + \int_{0}^{x_k} g_k(v) \, dv \right) = \frac{\lambda_k}{p_k(x_k)} G_k(x_k) ,
\]

(42)

(43)

with \( G_k(x_k) = \left( \pi_k^0 + \int_{0}^{x_k} g_k(v) \, dv \right) \) as in Eq. (11). Hence, (37) is valid with

\[
p_k(x_k) = \begin{cases} \frac{\lambda_k G_k(x_k)}{g_k(x_k)} & \text{if} \quad x_k > 0 , \\ 0 & \text{if} \quad x_k = 0 , \end{cases}
\]

(44)

\[
\pi_k(d x_k) = [\pi_k^0(\delta(x_k) + e^{-\zeta_k x_k} g_k(x_k))] \, d x_k .
\]

(45)

With this substitution, we obtain an equivalent formulation for the optimization problem in Eqs. (40a)-(40d) as below

\[
\sup_{\{ p_k^0(x_k) \}_{k=1}^{M}} \int_{0}^{T} r \left( \sum_{k=1}^{M} p_k(x_k) \right) \prod_{k=1}^{M} \pi_k(d x_k) ,
\]

(46a)

(46b)

(46c)

(46d)

where \( p_k(x_k) \) and \( \pi_k(d x_k) \) are according to (44) and (45).

Through the formulation in Eqs. (46a)-(46d), we can show that the throughput maximization problem in Eqs. (40a)-(40d) is concave with respect to each coordinate over a convex feasible set. In particular, since the transformation between \( f_k(x_k) \) and \( g_k(x_k) \) is linear, the concavity of (40a)-(40d) can be shown equivalently by proving the concavity of the formulation in Eqs. (46a)-(46d). To this end, suppose that \( \{ (\pi_k^{0.1}(x_k), \pi_k^{1.1}(x_k) \} \}_{k=1}^{M} \) and \( \{ (\pi_k^{0.2}(x_k), \pi_k^{1.2}(x_k) \} \}_{k=1}^{M} \) are two arbitrary sets of optimization parameters belonging to the feasible region defined in Eqs. (46b)-(46d). Then for all \( \alpha \in [0, 1] \) and \( \tilde{\alpha} \overset{\text{def}}{=} (1 - \alpha) \), it readily follows that \( \{ (\pi_k^{0.0}(x_k), \pi_k^{1.0}(x_k) \} \}_{k=1}^{M} \) also satisfies (46b)-(46d), where \( \pi_k^{0.0} = \alpha \pi_k^{0.1} + \tilde{\alpha} \pi_k^{0.2} \) and \( g_k^{0.0}(x_k) = \alpha g_k(x_k) + \tilde{\alpha} g_k(x_k) \) are the convex combinations of the densities and atoms, respectively. This proves the convexity of the feasible region (46b)-(46d).

Proposition 2. (Coordinate-wise Convexity) Let \( \bar{R} \eta, \bar{R} \eta_j \) and \( \bar{R} \eta_j \) be the utility functions corresponding to \( \{ (\pi_k^{0.0}(x_k), \pi_k^{1.0}(x_k) \} \}_{k=1}^{M} \), \( \{ (\pi_k^{0.1}(x_k), \pi_k^{1.1}(x_k) \} \}_{k=1}^{M} \), and \( \{ (\pi_k^{0.2}(x_k), \pi_k^{1.2}(x_k) \} \}_{k=1}^{M} \) respectively, such that

\[
\begin{cases} (\pi_k^{0.0}(x_k), g_k^{0.0}(x_k)) = \alpha(\pi_k^{0.1}(x_k), g_k^{1.0}(x_k)) + \tilde{\alpha}(\pi_k^{0.2}(x_k), g_k^{1.2}(x_k)) , \quad k = j , \\ (\pi_k^{0.0}(x_k), g_k^{0.0}(x_k)) = \alpha(\pi_k^{0.1}(x_k), g_k^{1.0}(x_k)) + \tilde{\alpha}(\pi_k^{0.2}(x_k), g_k^{1.2}(x_k)) , \quad k \neq j . \end{cases}
\]

Then,

\[
\bar{R} \eta_j \geq \alpha \bar{R} \eta_j + \tilde{\alpha} \bar{R} \eta_j .
\]

(47)

Proof. The proof is relegated to Appendix B.
Now, define an ensemble of perturbation functions, \( \{ \psi_k \}_{k=1}^M \), such that
\[
\int_{0^+}^{L_k} \psi_k(v) \, dv = 0 \quad \text{(48)}
\]
and the \( \psi_k \) are continuous and bounded on their domain \((0, L_k)\) and \( \psi_k(0) = 0 \). For sufficiently small \( \varepsilon_k > 0, k \in [M] \), it thus follows that \( g_k^\varepsilon(x_k) = \varphi_k(x_k) + \varepsilon_k \psi_k(x_k) \) satisfies (46b)-(46d) with the same atoms \( \{ \varphi_k \}_{k=1}^M \) and thus lies inside the feasibility region. Then, with \( \varepsilon = (\varepsilon_1, \varepsilon_2, \ldots, \varepsilon_M) \), it must be true for a global solution that
\[
\hat{R}^\varepsilon \leq \hat{R}, \quad \text{(50)}
\]
where
\[
\hat{R}^\varepsilon = \int_{\mathcal{A}} r \left( \sum_{k=1}^M p_k^\varepsilon(x_k) \right) \prod_{k=1}^M \pi_k^{\varepsilon_k}(dx_k) \quad \text{(51)}
\]
and
\[
\pi_k^{\varepsilon_k}(x_k) \overset{\text{def}}{=} \pi_k^0 \delta(x_k) + e^{-\varepsilon_k x_k} g_k(x_k) + \varepsilon_k e^{-\varepsilon_k x_k} \psi_k(x_k) \, dx_k
\]
and \( p_k^\varepsilon(x_k) \) is calculated from Eq. (44) to be,
\[
p_k^\varepsilon(x_k) = \begin{cases} 
\lambda_k G_k(x_k) + \varepsilon_k \Psi_k(x_k) & x_k > 0 \\
0 & x_k = 0,
\end{cases} \quad \text{(53)}
\]
with,
\[
\Psi_k(x_k) \overset{\text{def}}{=} \int_0^{x_k} \psi_k(v) \, dv. \quad \text{(54)}
\]
For the moment, we assume that only the \( j \)th coordinate is perturbed; that is \( \varepsilon_k = 0, \forall k \neq j \). Expanding the right hand side of (51) to first order then results in
\[
\hat{R}^\varepsilon = \int_{\mathcal{A}} \left[ r \left( \sum_{k=1}^M p_k(x_k) \right) + \varepsilon_j \frac{\partial r \left( \sum_{k=1}^M p_k(x_k) \right)}{\partial p_j(x_j)} \frac{d p_j^\varepsilon(x_j)}{dx_j} \right] \\
\times \left\{ \pi_j(dx_j) + \varepsilon_j e^{-\varepsilon_j x_j} \psi_j(x_j) \, dx_j \right\} \prod_{k \in [M]-j} \pi_k(dx_k)
\]
\[
= \hat{R} + \varepsilon_j \int_{\mathcal{A}} \left[ r \left( \sum_{k=1}^M p_k(x_k) \right) e^{\varepsilon_j x_j} \psi_j(x_j) \right] \, dx_j \prod_{k \in [M]-j} \pi_k(dx_k) + O(\varepsilon^2) \quad \text{(55)}
\]
On the other hand, we note that
\[
\frac{d p_j^\varepsilon(x_j)}{dx_j} \bigg|_{x_j=0} = 0 \quad \text{(56)}
\]
since \( p_j^\varepsilon(x_j) = 0 \) from (53). Therefore,
\[
\int_{\mathcal{A}} \left[ r \left( \sum_{k=1}^M p_k(x_k) \right) \frac{d p_j^\varepsilon(x_j)}{dx_j} \right] \, dx_j \bigg|_{x_j=0} = 0, \quad \text{(57)}
\]
and thus from Eq. (55) we obtain (58) on the next page. This expansion, accompanied with inequality (50) establishes (59) as a necessary condition for a locally (and thus globally) optimal solution, where we have neglected the second order term \( O(\varepsilon^2) \). Now with slight abuse of notation, let
\[
\mathbb{E}_j \left[ \left( \sum_{k=1}^M p_k(x_k) \right) \prod_{k \in [M]-j} \pi_k(dx_k) \right] \overset{\text{def}}{=} \int_{\mathcal{A}} r \left( \sum_{k=1}^M p_k(x_k) \right) \prod_{k \in [M]-j} \pi_k(dx_k)
\]
denote the expectation over all the coordinates except the \( j \)th coordinate. Then (59) can be restated as
\[
\int_0^{L_j} \left[ \frac{\partial \mathbb{E}_j \left[ r \left( \sum_{k=1}^M p_k(x_k) \right) \right]}{\partial p_j(x_j)} \right] d x_j \bigg|_{x_j=0} = e^{-\varepsilon_j x_j} g_j(x_j) + \mathbb{E}_j \left[ r \left( \sum_{k=1}^M p_k(x_k) \right) e^{-\varepsilon_j x_j} \psi_j(x_j) \right] d x_j = 0, \quad \text{(69)}
\]
where we used the fact that
\[
\mathbb{E}_j \left[ \frac{\partial r \left( \sum_{i=1}^n p_i(z_i) \right)}{\partial p_j(z_j)} \right] = \frac{\partial}{\partial p_j(z_j)} \mathbb{E}_j \left[ r \left( \sum_{i=1}^n p_i(z_i) \right) \right].
\]
On the other hand, from Eq. (69), i.e.,
\[
\int_0^{L_j} \left( \lambda_j e^{-\varepsilon_j x_j} \frac{\partial \mathbb{E}_j \left[ r \left( \sum_{k=1}^M p_k(x_k) \right) \right]}{\partial p_j(x_j)} \right) \psi_j(x_j) \, dx_j
\]
Remark

It is easy to verify that

\begin{equation}
\int \left[ e^{-\zeta x_j} \frac{\partial}{\partial p_j(x_j)} \right] \psi_j(x_j) \, dx_j
= \left( e^{-\zeta x_j} \frac{\partial}{\partial p_j(x_j)} \right) \psi_j(x_j) + e^{-\zeta x_j} \sum_{k=1}^{M} \psi_j(x_j) \right] \psi_j(x_j) \, dx_j = 0.
\end{equation}

(74)

holds only if the term inside the parentheses is in the form of $K_j \exp(-\zeta x_j)$ for some constant $K_j$. Thus, as a necessary condition, we have

\begin{equation}
p_j(x_j) p'_j(x_j) \frac{\partial^2 \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} p_k(x_k) \right) \right]}{\partial p_j(x_j)^2} + (\lambda_j - \zeta_j p_j(x_j))
\end{equation}

\begin{equation}
\times \frac{\partial \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} p_k(x_k) \right) \right]}{\partial p_j(x_j)} + \zeta_j \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} p_k(x_k) \right) \right] + K_j = 0.
\end{equation}

(79)

Remark 6. Rewriting Equation (79) as

\begin{equation}
p'_j(x_j) = -p_j(x_j) \frac{\partial^2 \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} p_k(x_k) \right) \right]}{\partial p_j(x_j)^2} \left( \lambda_j - \zeta_j p_j(x_j) \right)
\end{equation}

\begin{equation}
\times \frac{\partial \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} p_k(x_k) \right) \right]}{\partial p_j(x_j)} + \zeta_j \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} p_k(x_k) \right) \right] + K_j,
\end{equation}

(80)

it is easy to verify that $K_j$ provides a degree of freedom to set the initial slope $p'_j(x_j)|_{x_j=0}$ of the power policy $p_j(x_j)$. Now since the choice of $j^{th}$ coordinate was arbitrary, (79) holds for all $j \in [M]$. Accordingly, we obtain a system of coupled PIEDs over $M$ coordinates with $2M$ degree of freedom where the integration is implicit in the notation of $\mathbb{E}_j \left[ \cdot \right]$ (ref. Eq. (68)). In the following, we consider solutions in the infinite and finite battery cases.

1) Infinite Storage Battery: Motivated by the converse result for the infinite storage battery, we consider a set of admissible policies as below

\begin{equation}
\bar{p}_k(x_k) = \begin{cases} 
(\lambda_k/\zeta_k) + \varrho & x_k > 0 \\
0 & x_k = 0,
\end{cases}
\end{equation}

(81)

where the excess power $\varrho > 0$ is added to ensure the positive recurrence of the process. In the limit, as $\varrho \to 0$, the suggested policies in Eq. (81) satisfy (79) for all $j$, provided

\begin{equation}
K_j = -\zeta_j \mathbb{E}_j \left[ r \left( \sum_{k=1}^{M} \frac{\lambda_k}{\zeta_k} \right) \right].
\end{equation}

(82)

The average transmission power of (81) can then be evaluated as

\begin{equation}
\mathbb{E}[\bar{p}_k(x_k)] = \frac{\pi_k^0 \bar{p}_k(0) + (\lambda_k/\zeta_k + \varrho)}{\int_{0}^{\infty} f_k(x_k) \, dx_k}
\end{equation}

(83)

\begin{equation}
= ((\lambda_k/\zeta_k + \varrho)(1 - \pi_k^0).
\end{equation}

(84)

On the other hand, in light of Eq. (34) we have $\mathbb{E}[\bar{p}_k(x_k)] = \lambda_k/\zeta_k$. As a result, a transmission node that exploits $\bar{p}_k(x_k)$ as transmission power policy has the following probability mass at zero

\begin{equation}
\pi_k^0 = \frac{\varrho}{(\lambda_k/\zeta_k + \varrho)}.
\end{equation}

(85)
\[ \frac{\lambda_k}{\zeta_k} + \varrho \]

Moreover, associated with \( \bar{\rho}_k(x_k) \), the mean square deviation of transmission power is given by
\[ \sigma^2(\bar{\rho}_k(X_k)) = \frac{\lambda_k}{\zeta_k} \sigma_k^2 + \int_{0^+}^{\infty} (\bar{\rho}_k(x_k) - \frac{\lambda_k}{\zeta_k} f_k(x_k)) \, dx_k \]
\[ = (\lambda_k/\zeta_k) \varrho. \]  
(86)

For these power transmission strategies, we also have
\[ \hat{R} \geq r \left( \sum_{k=1}^{M} \frac{\lambda_k}{\zeta_k} \right) \]  
(87)

where the inequality follows from neglecting situations in which a strict subset of nodes are transmitting and the rest are silent due to battery exhaustion. As \( \varrho \downarrow 0 \), the upper bound (35) and the lower bound (87) coincide with each other. The total average throughput is given asymptotically by
\[ \hat{R} = r \left( \sum_{k=1}^{M} \frac{\lambda_k}{\zeta_k} \right). \]  
(88)

Thus, near optimal performance of the energy harvesting system can be achieved when \( \varrho \downarrow 0 \), and the behaviour of a classical communication systems (in the sense of using a constant power supply without interruption) can be closely approximated at the same time.

2) Finite Storage Battery: In contrast to the case of batteries with infinite capacity, the system of equations in Eq. (79) doesn’t appear to admit a closed form expression for power policies when the storage capacity is finite. The remaining option is thus to solve (79) numerically. However, the complexity in dealing with such systems is that the equations are not independent, but coupled. Alternatively, if all but one of the \( p_k(\cdot) \) are known, the remaining one can be obtained by solving a first order non-linear ODE in terms of the corresponding coordinate, using (79). First from (43) and since \( G_k(x_k) = g_k(x_k) \) we obtain by integration
\[ \ln G_k(x_k) - \ln G_k(0) = \int_{0^+}^{\infty} \frac{\lambda_k}{p_k(v_k)} \, dv. \]  
(90)

Therefore,
\[ G_k(x_k) = \pi_k^0 \exp \left( \int_{0^+}^{\infty} \frac{\lambda_k}{p_k(v)} \, dv \right). \]  
(91)

By differentiating both sides with respect to \( x_k \), we obtain
\[ g_k(x_k) = \frac{\pi_k^0 \lambda_k}{p_k(x_k)} \exp \left( \int_{0^+}^{\infty} \frac{\lambda_k}{p_k(v)} \, dv \right), \]  
(92)

or equivalently,
\[ f_k(x_k) = \pi_k^0 e^{-\lambda_k x_k} \exp \left( \int_{0^+}^{\infty} \frac{\lambda_k}{p_k(v)} \, dv \right). \]  
(93)

Also, due to the normalization condition
\[ \pi_k^0 + \int_{0^+}^{L_k} f_k(x_k) \, dx_k = 1, \]  
(94)

we have
\[ \pi_k^0 = \left[ 1 + \int_{0^+}^{L_k} e^{-\lambda_k x_k} \exp \left( \int_{0^+}^{\infty} \frac{\lambda_k}{p_k(v)} \, dv \right) \right]^{-1}, \]  
(95)

which can simply be derived via substituting (93) in Eq. (94) and solving for \( \pi_k^0 \).

With the help of (93) and (95), we propose an iterative method, outlined as Algorithm 1, that computes a solution for (79). Also, the convergence analysis of Algorithm 1 follows the following three steps:

(i) at each iteration of Algorithm 1, the utility function (37) is non-decreasing,

(ii) the utility function in Eq. (37) is bounded above,

(iii) the utility (37) thus converges if Algorithm 1 is allowed to iterate indefinitely (i.e. no termination constraint).

Specifically, in the first step, we denote the utility \( \hat{R} \) as an explicit function of the power policies in Algorithm 1, e.g., \( \hat{R}(p_1^{(0)}(x_1), p_2^{(0)}(x_2), \ldots, p_M^{(0)}(x_M)) \) is the initial utility. After the \( N \)th full iteration of steps 5-10 (outer loop) of Algorithm 1, in the \( j \)th iteration of 6-10 (inner loop), we then obtain Eq. (96) below. Therefore,
\[ \hat{R}(p_1^{(N+1)}(x_1), \ldots, p_j^{(N)}(x_j), \ldots, p_M^{(N)}(x_M)) \leq \hat{R}(p_1^{(N+1)}(x_1), \ldots, p_j^{(N+1)}(x_j), \ldots, p_M^{(N)}(x_M)). \]  
(97)

In addition, since the objective function is upper bounded by (28), we further have
\[ \hat{R}^{\sup} \stackrel{\text{def}}{=} \sup_{\{ \bar{p}_k(x_k) \}} \hat{R}(p_1(x_1), p_2(x_2), \ldots, p_M(x_M)) \leq \left( \sum_{k=1}^{M} \frac{\lambda_k}{\zeta_k} \left( 1 - \exp(-\zeta_k L_k) \right) \right). \]  
(98)

Concluding from (97) and (99), the sequence \( \hat{R}(\{p_k^{(N)}(x_k)\}_{N=0}^{\infty}) \) must converge in the limit as \( N \rightarrow \infty \).

Now, we concentrate on two important degenerate cases of our problem that can substantially reduce the computational burden of solving the PIEDs. In the first scenario, suppose that all the transmission nodes scavenge energy in the same manner. By this statement, we mean that the statistical parameters of all the energy harvesters are identical, i.e., \( \lambda_k = \lambda \) and \( \zeta_k = \zeta \)

\[ p_j^{(N+1)}(x_j) = \arg \max_{\xi} \hat{R}(p_1^{(N+1)}(x_1), \ldots, p_j^{(N+1)}(x_j), p_{j+1}^{(N+1)}(x_{j+1}), \ldots, p_M^{(N)}(x_M)). \]  
(96)
for all \( k = 1, 2, \ldots, M \). In the symmetric case, further assume that the batteries have identical capacities (\( L_k = L \)) and all transmitters employ the same power policy (\( p_k(x_k) = p(x_k) \)). Then, Equation (79) reduces to

\[
 p(x_j)p'(x_j) \frac{\partial^2 E_j[r(\sum_{k=1}^M p(x_k))] + (\lambda - \zeta p(x_j))}{\partial^2 p(x_j)} + (\lambda - \zeta p(x_j)) \\
 \times \partial E_j[r(\sum_{k=1}^M p(x_k))] + \zeta E_j[r(\sum_{k=1}^M p(x_k))] + K = 0,
\]

(100)

where \( j \) is arbitrary and chosen from \([M] \), and the operator \( E_j[\cdot] \) now simplifies as

\[
 E_j \left[ r(\sum_{k=1}^M p(x_k)) \right] = \int_{A_j} r(\sum_{k=1}^M p(x_k)) \prod_{k \in [M] \setminus j} \pi(dx_k).
\]

(101)

If we rearrange the terms in Equation (100), we have that

\[
 p(x_j) = F(p(x_j)),
\]

(102)

where the mapping \( F(\cdot) : C^1(0, L) \to C^1(0, L) \) is given by

\[
 F(p(x_j)) =
 p(0^+) - \int_0^{x_j} \left[ K_j + \zeta E_j[r(\sum_{k=1}^M p(v_k))] + (\lambda - \zeta p(v_j)) \right] \\
 \times \partial E_j[r(\sum_{k=1}^M p(v_k))] \left[ p(v_j) \frac{\partial^2 E_j[r(\sum_{k=1}^M p(v_k))]}{\partial^2 p(v_j)} \right]^{-1} dv_j.
\]

(103)

As a result, it follows that the desired \( p(x_j) \) is a fixed point of \( F(\cdot) \). This then suggests an alternative algorithm for this special case (see Algorithm 2).

Now in the second scenario, consider that there is only one transmitter in the communication system (i.e., a point-to-point setup). We thus have a simplified formulation as a necessary condition here, i.e.,

\[
 p(x_j)p'(x_j) \frac{d^2 r(p(x_j))}{dp^2(x_j)} + (\lambda - \zeta p(x_j)) \frac{dr(p(x_j))}{dp(x_j)} + r(x_j) + K = 0.
\]

(104)

As argued in [1], this is a second order, non-linear, autonomous ODE that can be solved numerically by employing linear multistep methods (e.g. Runge-Kutta or Adams-Bashforth). The next lemma demonstrates some properties of solutions to this ODE.

**Lemma 3.** Suppose \( K > -\zeta r(\lambda/\zeta) \) in Eq. (104), then for the Shannon rate function

(i) any solution \( p(x) \) is a strictly increasing function of \( x \) for \( x \geq 0 \), and \( p(x) \to \infty \) as \( x \to \infty \),

(ii) \( p(x) \) grows doubly exponentially fast as \( x \to \infty \).

\( \square \)

**Proof.** Solving (104) for \( p'(x) \), we have

\[
 p'(x) = \frac{(\lambda - \zeta p(x)) r'(p(x)) + \zeta r(p(x)) + K}{-p(x) r''(p(x))}.
\]

(105)

From concavity of the rate function as well as the first constraint on admissible power policies we have \( r''(p(x)) < 0 \) and \( p(x) \geq 0 \), respectively. Therefore, the denominator is always positive and \( p'(x) > 0 \) for all \( x \geq 0 \) iff

\[
 K > -\left[(\lambda - \zeta p(x)) r'(p(x)) + \zeta r(p(x))\right], \quad \forall x \geq 0.
\]

(106)

Moreover, it can be verified that

\[
 \frac{d}{dp} \left[ (\lambda - \zeta p(x)) r'(p(x)) + \zeta r(p(x)) \right] = (\lambda - \zeta p(x)) r''(p(x)).
\]

Hence, \( p(x) = \lambda/\zeta \) is a global maxima for the right hand side of (106). Replacing \( p(x) = \lambda/\zeta \) into (106), the numerator of (105) is then lower bounded by

\[
 K + \zeta r(\lambda/\zeta) > 0.
\]

Furthermore, since \( r(x) = \frac{1}{2} \log_2 \left( 1 + \frac{x}{N_0} \right) \), we upper bound the denominator by

\[
 -p(x) r''(p(x)) = \frac{1/N_0}{2 \ln 2} \frac{p(x)/N_0}{(1 + p(x)/N_0)^2} \leq \frac{1/N_0}{8 \ln 2}
\]

and thus \( p(x) \to +\infty \).

To prove the second part of Lemma 3, consider the substitution

\[
 p(x)/N_0 = \exp(S(x)),
\]

(107)

where \( S(x) \) increases since \( p(x) \) increases. Then we have

\[
 S(x) \to +\infty, \quad \text{as} \quad x \to \infty.
\]

(108)

Consequently, for the Shannon rate function we obtain

\[
 \frac{dr(p(x))}{dp(x)} = \frac{1}{2 \ln 2 \left( 1 + p(x)/N_0 \right)} \approx \frac{1/N_0}{2 \ln 2} \exp\left(-S(x)\right),
\]

(109)

and

\[
 \frac{d^2 r(p(x))}{dp^2(x)} = -\frac{1}{2 \ln 2 \left( 1 + p(x)/N_0 \right)^2} \approx -\frac{1/N_0^2}{2 \ln 2} \exp(-2S(x)),
\]

(110)

and

\[
 r(p(x)) = \frac{1}{2} \log_2 \left( 1 + p(x)/N_0 \right) \approx \frac{1}{2 \ln 2} S(x).
\]

(111)

Replacing (109)-(111) in Eq. (104) yields that for \( x \to \infty \)

\[
 -S'(x) + (\lambda - \zeta N_0 e^{S(x)}) (e^{-S(x)/N_0}) + \zeta S(x) + (K/2 \ln 2) = 0,
\]

(112)

As \( x \to \infty \) and due to (108), Equation (112) reduces to

\[
 \zeta S(x) = S'(x),
\]

(113)

which has the following solution

\[
 S(x) = A \exp(\zeta x),
\]

(114)

for some constant \( A \), and thus

\[
 p(x) = O(\exp(e^{\zeta x})), \quad \text{as} \quad x \to \infty.
\]

(115)
Remark 7. Due to Lemma 1, it is easy to verify that when \( K \leq -\frac{\lambda}{\zeta_a} \), the property of (115) does not hold in general. In fact, for sufficiently large negative \( K \), solutions of Eq. (104) are decreasing power policies. However, we conjecture that all such power policies are suboptimal as they fail to control the overflow in the battery. A more detailed discussion will be presented in the following section.

V. NUMERICAL EXPERIMENTS

We now study a multiple access communication system consisting of two nodes (\( M = 2 \)) with \( \lambda_1 = \lambda_2 = \lambda = 1 \) and \( \zeta_1 = \zeta_2 = \zeta = 1 \). Because of the symmetry of the MAC, the achievable power policies for this setting are obtained through Algorithm 2. However, to implement Algorithm 2 according to steps 1-6, one is obliged to search for optimized values of \( p(0^+) \) and \( K \) at each iteration. To ease this process and in what follows, Algorithm 2 is modified in a way that once the values of \( p(0^+) \) and \( K \) are initialized, the same values are used at each iteration step.

With this modification, Fig. 1(a) then shows the designed power policy as a function of the remaining charge in the battery with initial conditions \( p(0^+) = 0.1 \) and \( K = 0 \) and initializing function \( p^0(x_k) = x_k + p(0^+) \), \( 0 < x_k \leq L_k \). After \( N = 10 \) iterations, the power policy has converged to a solution of (100). It can also be seen from Fig. 1(a) that as the remaining charge in the battery increases, the transmission power also increases rapidly. Supported by part (b) of Lemma 3, we further conjecture that this increase is in fact doubly exponential in \( x \). Indeed, when the occupied charge of the battery becomes large, the chance of overflow due to new energy arrivals increases as well. In this regard, an optimal power policy is one which consumes the battery charge fast enough such that the occurrence of overflow is traded-off against the potential suboptimality of employing a large instantaneous transmission power (see Section IV). On the other hand, for sufficiently large negative \( K \), solutions of Eq. (100) are non-increasing (see Remark 7 for the point-to-point case) and they thus fail to manage battery overflow. The numerical results have further verified that for non-increasing power policies, the achieved sum-throughput is strictly less than for increasing ones. As a result, here we only consider increasing power policies.

Corresponding to the designed power policy in Figure 1(a), Figure 1(b) shows the absolutely continuous part (density) of the probability measure in Eq. (7). In this case, consistent with our earlier observation for the power policy, the density function also falls off quickly. In terms of ergodicity, this is basically an assertion of the fact that the system spends little time with large stored charge in the battery.

Using Algorithm 2, we have computed the achievable rates provided in Table I and Table II, where the termination criterion is taken to be

\[
\theta = \frac{r \left( \sum_{k=1}^{M} P^{N+1}(x_k) \right) - r \left( \sum_{k=1}^{M} P^N(x_k) \right)}{r \left( \sum_{k=1}^{M} P^N(x_k) \right)} < 1\%,
\]

i.e., the iteration stops whenever the increase in rate is less than one percent. With this precision, Table I shows the sum throughput for several choices of \( p(0^+) \) and fixed \( K = 0 \). The upper bound is also computed from (28) and denoted by \( R_{upper} \) in the table.

Based on a comparison between the upper and lower limits on the rate function, it is immediate that the choice of \( p(0^+) = 0.001 \) results the best performance of the designed power policy. For the same choices of \( p(0^+) \) as in Table I and \( K = 0 \), Figure 2 shows the power policy solutions. Except for the case of \( p(0^+) = 1 \), all the power policy solutions adopt a small transmission power when the battery charge is small. Along the same lines, Table II shows the upper and lower limits on the average throughput for fixed \( p(0^+) = 0.001 \) and variable \( K \). We have particularly provided the best value of \( K \) up to precision 0.01 as well as the corresponding achievable rates. The best achievable rate, as a percentage of the upper bound, is also evaluated.

Finally, to show the robustness of the iterative algorithm to the initialing function, a different choice of \( p_k^0(x_k) \) is studied in Figure 3. Therein, we particularly have selected \( p_k^0(x_k) = p_{k-1}^0 + \sqrt{x_k} \) for purpose of initialization in Algorithm 2 while the rest of the parameters are the same as in Figure 1(a). Evidently, the power policy converges to an identical function as one depicted in Figure 1(a). Similarly, the same convergence was observed when \( p_k^0(x_k) = p_k^0 + \sqrt{x_k} \), \( 0 < x_k \leq 3 \). In this respect, the proposed algorithm appears to be insensitive to the choice of the initial power policy.

VI. CONCLUSION

We have considered continuous-time power policies for a multiple access communication system where each node is capable of harvesting energy. First we modelled the battery as a compound Poisson dam, where the remaining charge in the battery modulates the transmission power. We then analysed this storage dam model in the ergodic case. In particular, we characterized an upper bound on the maximum sum-rate as a function of the energy arrivals distributions and the capacity of the batteries. For batteries with infinite capacity, we proved that any rate close to this bound is achievable by a set of constant power policies that result in stable battery behavior. For
batteries with limited capacity, we showed that optimal power policies can be derived by solving a system of simultaneous partial integro-differential equations. To solve these equations, we developed an iterative algorithm based on the Gauss-Seidel approach. We next derived a fixed point algorithm for the symmetric MAC case where the multiple access nodes have identical energy harvesting statistics. Furthermore, the convergence of the utility function that results from the proposed algorithms was established. Numerical results show that for $L = 3$, the achievable scheme provides throughput up to 86.6% of the upper bound.

Potential future work includes extending the study to the case where each transmitter has a data buffer. This could model scenarios in which a sensor monitors a physical quantity (e.g., temperature), and then stores the data in a buffer for eventual transmission once enough energy has been harvested.

**APPENDIX A**

In the following, we show that for every power policy with memory, $p_k^*(X_k(u); u \leq t)$, there exist a memoryless counterpart $p_k(X_k(t))$ that attains the same or better sum-throughput performance. In particular, let $(\Omega, \mathcal{F}, \mathbb{P})$ be a complete probability space and $X_k(t; \omega)$, $-\infty < t < \infty$ be a stationary and ergodic stochastic process defined on this probability space and whose evolution for $t \geq 0$ is given by

$$X_k^*(t; \omega) = X_k^*(0; \omega) + E_k^{\text{ln}}\left((0, t]; \omega\right) - \int_0^t p_k^*(X_k^*(u; \omega); u \leq s)\,ds. \tag{116}$$

In conjunction with the process $X_k^*(t; \omega)$, we then define the following empirical CDFs,

$$\tilde{F}_k^*(\rho_k, x_k; \omega) \equiv \lim_{T \to \infty} \frac{1}{T} \int_0^T 1(p_k(X_k^*(u; \omega); u \leq s) \leq \rho_k) 1(X_k^*(s; \omega) \leq x_k)\,ds$$

$$\tilde{\pi}_k^*(x_k; \omega) \equiv \lim_{T \to \infty} \frac{1}{T} \int_0^T 1(X_k^*(s; \omega) \leq x_k)\,ds = \lim_{\rho_k \to \infty} \tilde{F}_k^*(\rho_k, x_k; \omega). \tag{119}$$

Now since $X_k^*(t)$ is ergodic, $\tilde{F}_k^*(\rho_k, x_k; \omega)$ and $\tilde{\pi}_k^*(x_k; \omega)$ are well defined, and constant $\mathbb{P}$–almost surely on $\Omega$, i.e.,

$$\tilde{F}_k^*(\rho_k, x_k; \omega) \stackrel{\text{a.s.}}{=} F_k^*(\rho_k, x_k), \tag{120a}$$

$$\tilde{\pi}_k^*(x_k; \omega) \stackrel{\text{a.s.}}{=} \pi_k^*(x_k), \tag{120b}$$

For the functions $F_k^*(\rho_k, x_k)$, $\pi_k^*(x_k)$ in Eqs. (120) we define the conditional CDF $F_k^*(\rho_k | x_k)$ by

$$F_k^*(\rho_k, x_k) = \int_0^{x_k} F_k^*(\rho_k | s)\pi_k^*(ds). \tag{121}$$

Also, we define the memoryless power policy $p_k(x_k)$ as follows

$$p_k(x_k) \equiv \int_0^{\infty} \rho_k F_k^*(d\rho_k | x_k), \tag{122}$$

and a corresponding storage process $X_k(t)$ governed by

$$X_k(t) = X_k(0) + E_k^{\text{ln}}\left((0, t]; \omega\right) - \int_0^t p_k(X_k(s))\,ds. \tag{123}$$

Also, we denote the stationary measure of $X_k(t)$ by $\pi_k(x_k)$. Our objective now is to prove that the throughput using the storage process with memory $p_k^*(X_k^*(u; \omega), u \leq t)$ is no better than that of it’s memoryless counterpart $p_k(X_k(t))$, i.e.,

$$\hat{R}(\{p_k^*(X_k^*(s; \omega); s \leq t)\})_{k=1}^M \leq \hat{R}(\{p_k(X_k(t))\})_{k=1}^M, \tag{124}$$

almost surely. To show this result, we begin with the definition of the long term average throughput for the storage process in Eq. (116), i.e.,

$$\hat{R}\left(\{p_k^*(X_k^*(s; \omega); s \leq t)\}_{k=1}^M\right) \equiv \lim_{T \to \infty} \frac{1}{T} \int_0^T r\left(\sum_{k=1}^M p_k^*(X_k^*(s; w); s \leq t)\right)\,ds \tag{125}$$

$$\int_{\mathcal{A}} \int_{\mathcal{B}} r(X_k(s; w)) \prod_{k=1}^M \tilde{F}_k^*(d\rho_k, dx_k; w) \tag{126}$$

For clarity, we make the dependence on $\omega \in \Omega$ explicit.
we obtain corresponds to negative jumps in the sample.

\[ \int_0 T^0 \leq \int_A \left( \sum_{k=1}^M r \left( \sum_{k=1}^M \rho_k \prod_{t=1}^M F^*_k (d \rho_k | x_k) \right) \prod_{k=1}^M \pi^*_k (dx_k) \right) \]

\[ = \int_A \left( \sum_{k=1}^M F^*_k (d \rho_k | x_k) \right) \prod_{k=1}^M \pi^*_k (dx_k) \quad (129) \]

\[ \text{def} \int_A \left( \sum_{k=1}^M F^*_k (d \rho_k | x_k) \right) \prod_{k=1}^M \pi^*_k (dx_k), \quad (130) \]

where the last step follows from the definition of \( p_k(x_k) \) in Eq. (122). The remaining task is now to show that \( \pi^*_k (x_k) = \pi_k (x_k), \forall k \in [M] \). For this purpose, we define some notation in conjunction with an arbitrary, stationary, \( \mathcal{C} \)ad\( \mathcal{C} \)lag process \( Y(t) \) whose jumps (positive or negative) occur at time instants \( T^0, T^1, \cdots \). In particular, the right hand derivative of \( Y(t) \) is defined by

\[ Y^+(t) \overset{\text{def}}{=} \lim_{\varepsilon \downarrow 0} \frac{Y(t + \varepsilon) - Y(t)}{\varepsilon}. \quad (131) \]

In addition, define

\[ Y^-(t) \overset{\text{def}}{=} \lim_{\varepsilon \downarrow 0} Y(t - \varepsilon). \quad (132) \]

**Theorem 4.** (Rate Conservation Law) Let \( Y(t) \) be an ergodic, stationary, \( \mathcal{C} \)ad\( \mathcal{C} \)lag process. Then,

\[ f(y) \mathbb{E}^0 \left[ Y^+(t) \mid Y(t) = y \right] = \lambda^0 \mathbb{E}^0 \left[ 1_{\{Y(T^0 ) < y\}} 1_{\{Y(T^0 ) < y\}} - 1_{\{Y(T^0 ) < y\}} 1_{\{Y(T^0 ) > y\}} \right], \]

where \( f(y) \) is the probability density at \( y \), and \( \mathbb{E}^0 \) denotes the expectation with respect to the Palm probability distribution corresponding to the point process (with assumed intensity \( \lambda^0 \)) for the jumps. \( \square \)

**Proof.** The proof can be found in [23, p. 36]. \( \blacksquare \)

**Remark 8.** The term \( 1_{\{Y(T^0 ) < y\}} 1_{\{Y(T^0 ) < y\}} \) in the right hand side of Theorem 4 corresponds to negative jumps in the sample path while \( 1_{\{Y(T^0 ) < y\}} 1_{\{Y(T^0 ) > y\}} \) corresponds to positive jumps.

**Remark 9.** As the memoryless storage process in Eq. (123) only contains positive jumps,

\[ \mathbb{E}^0 \left[ 1_{\{X_k (T^0 ) > x_k \}} 1_{\{X_k (T^0 ) < x_k \}} \right] = 0, \]

where as defined in Section II-B, \( T_k^0, T_k^1, \cdots \) denote the energy arrival times for the \( k^{th} \) node. In this special case, we then have

\[ f_k(x_k) \mathbb{E}[X^*_k(t)|X_k(t) = x_k] = f_k(x_k) \mathbb{E}[\{p_k(X_k(t))|X_k(t) = x_k\}] - f_k(x_k) p_k(x_k). \quad (134) \]

For the right hand side of Theorem 4 we obtain

\[ \lambda^0 \mathbb{E}^0 \left[ -1_{\{X_k (T^0 ) < x_k \}} 1_{\{X_k (T^0 ) > x_k \}} \right] \overset{(a)}{=} \lambda_k \mathbb{E}^0 \left[ -1_{\{X_k (T^0 ) < x_k \}} 1_{\{X_k (T^0 ) > x_k \}} \right] \]

\[ = -\lambda_k \int_0^x \left( 1 - B_k(x_k - v_k) \right) \pi_k (dv_k) \quad (135) \]

\[ = -\lambda_k \left[ 1 - B_k(x_k) \right] \pi_k^0 + \int_0^x \left( 1 - B_k(x_k - v_k) \right) f(v_k) \, dv_k \quad (136) \]

Right continuous with left hand limit. Note that both the storage processes in Eqs. (116) and (123) are \( \mathcal{C} \)ad\( \mathcal{C} \)lag.
where (a) follows from the notion of Poisson Arrivals See Time Averages (PASTA) [23, Prop. 1.23] for Poisson energy arrival process with intensity $\lambda_0 = \lambda_k$. Equating (134) and (137) according to Theorem 4, we obtain

$$f_k(x_k)p_k(x_k) = \lambda_k \left[ (1 - B_k(x_k)) \pi_k^0 \right] + \int_{0^+}^{x_k} (1 - B_k(x_k - v_k)) f_k(v_k) dv_k,$$

which is the equilibrium condition in Eq. (9) with the density $f_k(x_k)$ and the atom $\pi_k^0$.

Returning to the storage process with memory in Eq. (116), now it also easy to see that

$$\mathbb{E}[X_k^+(t) | X_k(t) = x_k] = \int_0^\infty \rho_k F_k(d\rho_k | x_k)$$

which is simply the average rate of down crossing at level $x_k$ corresponding to the stationary distribution of $X_k^+(t)$ in Eq. (116). Using the argument in Remark 9 for the process in Eq. (116) results

$$f_k^+(x_k) \mathbb{E}[X_k^+(t)^+ | X_k^+(t) = x_k] = f_k^+(x_k)p_k(x_k)$$

$$= \lambda_k \left[ (1 - B_k(x_k)) \pi_k^0 \right] + \int_{0^+}^{x_k} (1 - B_k(x_k - v_k)) f_k^+(v_k) dv_k,$$

where $\pi_k^0$ and $f_k^+(x_k)$ are the atom and the continuous part (density) of the probability measure $\pi_k^+(x_k)$. Since from Theorem 1 the probability measure that solves (138) and (141) is unique,

$$\pi_k^+(x_k) = \pi_k(x_k), \quad \forall x_k.$$

Concluding from (130), we thus showed that

$$\hat{R}(\{p_k^i(X_k(s; w); s \leq t)\}_{k=1}^M) \leq \int_A \left( \sum_{k=1}^M \rho_k F_k(d\rho_k | x_k) \right) \pi_k^+(dx_k),$$

$$= \int_A \left( \sum_{k=1}^M \rho_k p_k(x_k) \right) \pi_k^+(dx_k),$$

$$\leq \hat{R}(\{p_k(X_k(t))\}_{k=1}^M).$$

Remark 10. We note that in the ergodic regime, the upcrossing rate as well as the drift component of the storage process in the finite battery case also obey the law stated in Theorem 4. Thus, we again obtain Eq. (141) as battery overflow does not change the upward and downward rates. Therefore, a similar proof can be used to show $X_k(t)$ is a sufficient statistic for optimal power policies in the storage model with a finite battery capacity in Eq. (6).

**APPENDIX B**

**Proof.** Since for all $k \neq j$,

$$\pi_k^{0,\alpha}(x_k) = \pi_k^{0,\beta}(x_k) = \pi_k^{0,\beta}(x_k),$$

we have that

$$\pi_k^+(dx_k) = \pi_k^+(dx_k) = \pi_k^+(dx_k).$$

Then,

$$\hat{R}_j^+ = \int_A r \left( \lambda_k \frac{g_j^k(x_k)}{g_j^k(x_k)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \pi_k^+(dx_k) \prod_{k \in [M]-j} \pi_k^+(dx_k),$$

$$= \mathbb{E}_j \left[ \int_0^{t\rightarrow} r \left( \lambda_k \frac{g_j^k(x_k)}{g_j^k(x_k)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \pi_k^+(dx_k) \right],$$

$$= \mathbb{E}_j \left[ \int_0^{t\rightarrow} r \left( \lambda_k \frac{g_j^k(x_k)}{g_j^k(x_k)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \pi_k^+(dx_k) \right]$$

$$\times \left[ \pi_j^{0,\alpha}\delta(x_j) + e^{-\gamma\varepsilon_j} g_j^0(x_j) \right] dx_j,$$

$$= \mathbb{E}_j \left[ \int_0^{t\rightarrow} r \left( \lambda_k \frac{g_j^k(x_k)}{g_j^k(x_k)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \pi_k^+(dx_k) \right]$$

$$+ \pi_j^{0,\alpha} \mathbb{E}_j \left[ \left( \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \right].$$

For the term inside the first expectation in Eq. (145), we proceed as (146)-(149) on the next page, where (148) can be verified via the lemma given in Appendix C and choosing

$$a_1 = \alpha e^{-\gamma\varepsilon_j} g_j^1(x_j), \quad a_2 = \alpha e^{-\gamma\varepsilon_j} g_j^2(x_j),$$

$$b_1 = \alpha e^{-\gamma\varepsilon_j} g_j^1(x_j), \quad b_2 = \alpha e^{-\gamma\varepsilon_j} g_j^2(x_j),$$

and

$$\gamma = \lambda_j, \quad \beta = \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)}.$$
\[
\int_0^{L_j} r \left( \lambda_j \frac{G_j^0(x_j)}{g_j^0(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) e^{-\zeta_j x_j} g_j^0(x_j) \, dx_j
\]
\[
= \int_0^{L_j} r \left( \lambda_j \frac{G_j^1(x_j)}{g_j^1(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \left( \alpha e^{-\zeta_j x_j} g_j^1(x_j) + \bar{\alpha} e^{-\bar{\zeta}_j x_j} g_j^2(x_j) \right) \, dx_j
\]
\[
= \int_0^{L_j} r \left( \lambda_j \frac{G_j^1(x_j)}{g_j^1(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) \left( \alpha e^{-\zeta_j x_j} g_j^1(x_j) + \bar{\alpha} e^{-\bar{\zeta}_j x_j} g_j^2(x_j) \right) \, dx_j
\]
\[
\geq \int_0^{L_j} r \left( \lambda_j \frac{G_j^1(x_j)}{g_j^1(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) e^{-\zeta_j x_j} g_j^1(x_j) \, dx_j
\]
\[
+ \bar{\alpha} \int_0^{L_j} r \left( \lambda_j \frac{G_j^2(x_j)}{g_j^2(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) e^{-\bar{\zeta}_j x_j} g_j^2(x_j) \, dx_j.
\]

\[
=E_j \left[ \int_0^{L_j} r \left( \lambda_j \frac{G_j^0(x_j)}{g_j^0(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) e^{-\zeta_j x_j} g_j^0(x_j) \, dx_j \right]
\]
\[
\geq \alpha E_j \left[ \int_0^{L_j} r \left( \lambda_j \frac{G_j^1(x_j)}{g_j^1(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) e^{-\zeta_j x_j} g_j^1(x_j) \, dx_j \right]
\]
\[
+ \bar{\alpha} E_j \left[ \int_0^{L_j} r \left( \lambda_j \frac{G_j^2(x_j)}{g_j^2(x_j)} + \sum_{k \in [M]-j} \lambda_k \frac{G_k(x_k)}{g_k(x_k)} \right) e^{-\bar{\zeta}_j x_j} g_j^2(x_j) \, dx_j \right].
\]

Furthermore, from concavity of \( V(x) \),
\[
b \sum_{k} (b_k/b)V(a_k/b_k) \leq b V \left( \sum_{k} \frac{b_k}{b} \times a_k/b_k \right)
\]
\[
= b V \left( \frac{a}{b} \right)
\]
\[
= \frac{\gamma}{\alpha} \left( b + \beta \right).
\]

Hence,
\[
\sum_{k} a_k r \left( \frac{b_k}{a_k} + \beta \right) \leq \frac{b}{\alpha} \left( b + \beta \right).
\]

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