Branching structure of genealogies in spatially growing populations and its implications for population genetics inference

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Abstract
Spatial models where growth is limited to the population edge have been instrumental to understanding the population dynamics and the clone size distribution in growing cellular populations, such as microbial colonies and avascular tumours. A complete characterization of the coalescence process generated by spatial growth is still lacking, limiting our ability to apply classic population genetics inference to spatially growing populations. Here, we start filling this gap by investigating the statistical properties of the cell lineages generated by the two dimensional Eden model, leveraging their physical analogy with directed polymers. Our analysis provides quantitative estimates for population measurements that can easily be assessed via sequencing, such as the average number of segregating sites and the clone size distribution of a subsample of the population. Our results not only reveal remarkable features of the genealogies generated during growth, but also highlight new properties that can be misinterpreted as signs of selection if non-spatial models are inappropriately applied.

Keywords: spatial growth, population genetics, branching process, KPZ universality class, site frequency spectrum, spatial coalescence

(Some figures may appear in colour only in the online journal)

1. Introduction

Spatial range expansions [1, 2] are ubiquitous in nature, from microbial biofilms [3, 4], developing tissues [5] avascular tumors [6–9] to invading species and infectious diseases [10, 11]. Many of these scenarios share the feature of being resources-limited [12–19], so that population growth occurs mainly as invasion of surrounding virgin territory [20–26].

When dispersal is local, these range expansions lead to a phenomenon called gene surfing, whereby pioneering individuals at the edge of the expansion have a higher chance to contribute to the next generation [27, 28]. As a result, an individual’s location can become a more important factor to reproductive success than its growth rate [29–32].

It has recently been shown that gene surfing leaves a characteristic signature in the mutational spectrum of the population, identified by an excess of high frequency mutations compared to the well-mixed expectation [4]. This observation becomes crucial when analyzing population sequencing results, as the same signature can be mistakenly interpreted as being a result of positive selection and lead to a mis-identification of driver
mutations (e.g., in cancer or drug resistance). Modeling the effects of spatial structure on genealogies in growing populations, and consequently on the diagnostic outputs of genome sequencing, could point to protocols that discern the two scenarios.

In prior efforts to link spatial growth with population structure, connections to non-equilibrium statistical mechanics have proven fruitful. Spatial population growth is a fundamentally out-of-equilibrium process driven by stochastic division and migration events at small scales that leave collective signatures at the scale of the population as a whole. Models of spatial populations harbor nonequilibrium statistical phenomena such as fixation into absorbing states [33], dynamic phase transitions and critical phenomena [34, 35], and manifestations of directed percolation [36]. As a result, statistical properties of the population patterns generated by the range expansion can be quantitatively linked to robust universal features of the corresponding nonequilibrium growth models [37]. A prime example is the connection between range expansions and the Kardar–Parisi–Zhang (KPZ) model of interface growth [38], which established scaling exponents for the shapes of clonal domains [20, 39, 40] and lineages [41–43] in expanding microbial populations. Such scaling rules provide a theoretical basis for predicting or interpreting population genetics quantities that can be measured using genome sequencing studies.

In the following, we systematically analyse the statistical properties of the genealogical tree generated by the Eden model, a lattice model that has successfully been used to investigate microbial colonies and tumour growth [4, 6, 44, 45], to determine the effects of spatial growth on three classic population genetics quantities: (i) time to the most recent common ancestor (MRCA), (ii) number of segregating sites in a sample and (iii) clone size distribution of a sample. We find that these quantities are completely determined by the growth properties of the populations, and that their key features can be captured by a deterministic tree structure defined completely by the exponents of the KPZ universality class. We finally discuss how recent advances in lineage tracing [14, 17, 46–50], as well as single cell sequencing can be combined with our model to reveal the presence of surface-limited growth and interpret the data accordingly.

2. Simulating spatial growth: the Eden model

The Eden model, first introduced in the seminal paper by Eden in 1961 [51], is widely used to mimic spatial growth processes where replication is limited to the front of the expansion, for example microbial colonies on rich media. Starting from an initial set of cells, placed at fixed points on a lattice, one cell with at least one empty neighbour is randomly chosen and replicated into one of the empty neighboring sites. This new cell can be seen as the descendant of the initially chosen cell and the process repeated to reach a final population size.

The growth process simulated by the model can be tracked to generate a genealogical tree that identifies the mother–daughter relationship of each individual (lattice site). The statistical properties of the emerging lineages have recently been investigated [43] and found to fall within the KPZ universality class.

This underlying growth process and the resulting lineages are sufficient to completely characterise the neutral genetic diversity of the population, since neutral mutations do not affect growth. The occurrence of neutral distinct mutations can be modeled as a Poisson stochastic process occurring on top of the identified lineages (infinite site model [52]). Using the statistical properties of the lineages, we can then characterise the corresponding coalescence process and estimate classic population genetics quantities. In what follows, we consider two scenarios: a linear front of constant width, which we compare to a Wright–Fisher model of constant population size; and a 2D radial expansion, which mimics colony growth. Details of the Eden model simulations are provided in appendix A.

3. Statistical properties of Eden model in a linear geometry

We will start our analysis with a linear front scenario (corridor) so that the front of the population exhibits a constant width. In what follows, we will always sample individuals from the very front of the population.

3.1. Number of segregating sites between two individuals

Single-cell sequencing enables genomic comparison (either whole-genome or targeted regions) among individual cells sampled from different locations in the population. The number of differences between the two (or more) genomes is a well-studied summary statistics in population genetics called number of segregating sites, S, whose distribution is known for well-mixed populations and even for simple models of structured populations (island model) [52].

In contrast to the well-mixed scenario, the spatial structure of our model naturally raises the question of how the number of segregating sites S depends on the relative location of the sampled individuals. Starting with two individuals sampled at a distance d from each other, then the probability $P_2(S|d)$ of observing S segregating sites is

$$P_2(S|d) = 2 \int P_1(S|T)P(T|d)dT,$$  \hspace{1cm} (1)

where $P_1(S|T)$ is the conditional probability of observing S segregating sites given that the time to the MRCA between the two individuals is T, $P(T|d)$ is the conditional probability of observing a time to the MRCA equal to T given that the two individuals are sampled a distance d apart. The factor 2 takes into account that mutations distinguishing the two individuals can occur on either branch leading to the MRCA. If we make the standard assumption that mutations follow a Poisson process, then $P_1(S|T) = \exp(-\mu T)(\mu T)^S/S!$, where $\mu$ is the mutation rate per replication and T is in units of replication events. In the Eden model (see appendix A), a replication event corresponds to the colonization of a neighboring lattice site, consequently both distances and times can be expressed in units of lattice sites.
The distribution $P(T|d)$ from linear simulations is shown in figure 1, with a power-law decay at large distances following $P(T|d) \sim T^{-1.64}$. The exponent is connected to one of the characteristic exponents describing the statistics of directed polymers in random media (DPRM) [43]. As $d$ increases, simulations deviate from the power-law expectation at large times due to the finite size of the simulations.

Because of the heavy tail of the distribution, the average time to the MRCA is often not of practical use, and the typical time to the MRCA, $\hat{T}$, is better suited as a metric of the characteristic behavior. This characteristic time has a scaling determined by the DPRM wandering exponent, $\hat{T} \sim d^{3/2}$ (figure 1(b)) [43].

Since the probability function $P(T|d)$ decays quickly upon moving away from the characteristic value $\hat{T}$, a simplified model of the tree structure can be built by replacing the distribution with a $\delta$-function peaked at $\hat{T}$. In addition, for large $\mu T$, we can also approximate $P_i(S|T)$ to a $\delta$-function peaked at the mean value $\mu T$, so the distribution of segregating sites scales similarly to $P(T|d)$, rescaled by a factor $2\mu$. In particular, the most likely number of segregating sites observed $\hat{S} \approx 2\mu \hat{T} \propto d^{3/2}$, following the KPZ expression (example 1). For comparison, in a Wright–Fisher model of constant population size $N$, the typical time to the MRCA would be $N$, independently on the physical distance between the two sampled individuals [52]. If we equate the population size $N$ to approximately the width $w$ of the corridor, this leads to a critical distance $d^* = N^{2/3} < w$, so that if individuals are sampled at distance larger than this, they should show more segregating sites than the well-mixed expectation, and vice versa.

### 3.2. Number of segregating sites in a subsample

If we consider a connected subsample $n < N$ at the front of the population, the total number of segregating sites is related to the total length of all the branches in the genealogical tree $T_{\text{tot}}(t)$. In the Wright–Fisher model this leads to the well-known average result

$$T_{\text{tot}} = \sum_{i=2}^n i T_i = 2N \log(n-1),$$

for an haploid population, where $T_i = N/ \binom{i}{2}$ represents the average time for the first coalescent event between two lineages of the possible $i$ [53]. Because after each coalescence event the number of lineages decreases by one, the total length of the tree is just given by the sum of the number of surviving lineages between subsequent coalescent events. Note that the time to the MRCA across $n$ individuals in a well-mixed population of size $N$ is $T_{\text{MRCA}}(n) = \sum_{i=2}^n T_i = 2N[1 - 1/(n - 1)]$.

For the Eden model, the expression depends on the relative position of the $n$ individuals. If we assume that they are positioned contiguously along the front, then $T_{\text{tot}}$ is the total length of the branches that lead to a corridor of width $n$ starting from the MRCA of the $n$ individuals (as in the inset of figure 2).

Then

$$T_{\text{tot}} = \int_1^{T_{\text{MRCA}}(n)} l(t) dt,$$

where $l(t)$ is the number of lineages at time $t$ measured backwards from the subsample and $T_{\text{MRCA}}(n)$ is the time to the MRCA of the whole sample. Because of the spatial constraints on the lineages, we have that $T_{\text{MRCA}}(n) \propto n^{3/2}$, since the $n$ individuals will be at most $n$ lattice sites apart (in reality there are more than $n$ individuals in a width $n$ since the front is rough, but we use this as first approximation).

The scaling of the typical time to the MRCA implies a particular scaling for the number of lineages with reverse time. We assume that all coalescence events happen exactly at the typical time $\hat{T}(d) \sim d^{3/2}$ associated with the separation $d$ of two individuals on the front. Upon advancing backwards in time from the front by an interval $t$, each contiguous segment of the boundary of size $\sim t^{2/3}$ will have coalesced to a single ancestor. As a result, the number of surviving lineages $l(t)$ falls as $t^{-2/3}$ with reverse time $t$, as long as the number of surviving lineages is large. To test this relation, we measured...
the number of lineages as a function of distance $s \propto t$ from the edge of the colony in Eden simulations within a linear corridor (see figure 2). Although the true genealogies display a more complex structure compared to our idealized model, we find that the observed behaviour in $l$ is consistent with a power-law decay $l \sim s^{-2/3} \sim t^{-2/3}$ until only a few lineages survive ($l \lesssim 1$).

Since the number of lineages is $l(t) \sim n/t^{2/3}$, then

$$T_{\text{tot}} \sim nT^{1/3}_{\text{MRCA}}(n) \propto n^{1/2}. \quad (4)$$

A remarkable byproduct of this is that the total tree size is proportional to the $T_{\text{MRCA}}(n)$ of the sample, unlike the case for well-mixed populations, and it does not depend on the total population size. This reflects the fact that the dominant contribution to the tree is given by the significantly longer oldest branches. Another interesting feature arising from the comparison with the Wright–Fisher model is that the total number of segregating sites increases much more quickly with sample size for the spatial model than for the well-mixed one.

### 3.3. Site frequency spectrum of subsamples

The statistical properties of the number of segregating sites in a sample $n$ of the edge of the population also determine the mutational spectrum, a commonly used genomic metric of the population structure. In the case of the Wright–Fisher infinite site model, the mutational spectrum, i.e., the number of mutations $m(j)$ carried by $j < N$ individuals, is given, on average, by $m(j) = 2\mu N/j$ [52].

The hierarchical length structure of the genealogical tree generated by the Eden model generates a very different mutational spectrum, since mutations can accumulate for a long time on long lineages before any later branching event occurs. Importantly, to understand the origin of this mutational spectrum, the topology of the tree (which branches coalesce with each other) is crucial, as identical $l(t)$ can generate very different mutational spectra. Furthermore, the measured spectrum will depend strongly upon the spatial distribution of the population samples—different sampling protocols might be sensitive to different features of the genealogical structure.

If samples are taken uniformly across the entire population, the site frequency spectrum is expected to follow a trend $m(j) \propto j^{-7/5}$ [4]. However, in many situations the outer edge of the population is more accessible for sampling. An analogous theoretical argument to the one used in [4] can be made if we restrict sampling to the edge of the population. Then, mutations carried by at least $j$ individuals have to occur before the coalescence time $T_{\text{MRCA}}(j)$ or, in other words, somewhere on the subtree between the ancestor among all $N$ individuals and the ancestor of the subsampled $j$ individuals. We will call this portion of the total tree $T^*(j)$. The topology of the tree then determines the expression for the site frequency spectrum.

For instance, if the tree is well-balanced so that coalescence events happen almost at the same time between pairs of lineages that are the same distance apart, then the number of mutations $m(j)$ carried by at least $j$ individuals is proportional to

$$T^*(j) = T_{\text{tot}}(N) - T_{\text{tot}}(j)N/j, \quad (5)$$

as there would be $N/j$ identical subtrees emerging from the corresponding $N/j$ lineages, each with $j$ leaves on average. This leads to

$$T^*(j) \approx N^{3/2} - j^{3/2}N/j = N^{3/2}(1 - x^{1/2}), \quad (6)$$

where $x = j/N$ is the frequency of the mutation. This expression leads to a clone size distribution $\Pi'(x)$ (probability that a mutation is carried by a proportion $x$ of the population) to be proportional to $x^{-1/2}$.

Eden model simulations show that, at least for small $x$ when finite size effects are limited, the scaling is $\Pi'(x) \propto x^{-2/3}$ (figure 3), clearly indicating that the tree, in this case, is not balanced, and the tail of the $T_{\text{MRCA}}$ distribution plays...
4. Radial expansion

While linear expansions are useful to understand the properties of the tree structure generated by the spatial growth process, radial expansions are more relevant to several real case scenarios (e.g., microbial colonies, avascular tumors). In this case, the population expands initially very rapidly due to an inflation effect related to curvature, which slows down as the radius grows [33].

From radial Eden model simulations, we find that the number of lineages as a function of the distance \( s \) from the edge of the colony follows the same \( s^{-2/3} \) power-law as in the corridor case, with a sharp drop close to the centre of the colony where the lineages spread star-like due to the rapid inflation process (figure 4).

To account for the inflation process, we normalize the number of lineages present at radius \( r \) by the circumference of the colony at the same radius to obtain a lineage density, and rescale the radius \( r \) by the final colony radius \( R \) (figure 5). We observe that consistently across colony sizes, the lineage density undergoes a transition at exactly a radius \( R/2 \), where \( R \) is the radius of the final colony, so that the density of surviving lineages initially decreases and then increases with \( r \). This non-monotonic behaviour reflects the tradeoff between the process of inflation, which pushes lineages apart preventing them from coalescing, and the stochastic wandering of the lineages, which over time makes them coalesce.

The collapse of the lineage density on a master curve independently of colony size suggests the presence of a universal tree that can describe the behaviour of genealogies generated in these two-dimensional spatial growth models. We propose it below.

4.1. Universal tree model

Our proposed model incorporates the statistics of lineage fluctuations imposed by the KPZ universality class, which the Eden model is known to belong to, as well as the spatial constraints on coalescence arising from the radial structure of the expansion (figure 6).

Building on our results for the corridor case, we hypothesize that the lineages of two cells which lie a spatial arc distance \( d \) apart at a radius \( r \gg d \) from the center of the colony will most likely coalesce at a certain distance \( h \) towards the center of the colony, so that

\[
h \sim d^{3/2}.
\]

We will now neglect rare stochastic events in the coalescence process of the colony’s lineages and devise a deterministic model where coalescence is controlled only by the typical coalescence height. The model is based on a binary tree with its branches’ lengths following the above relation (7). We will also assume that always exactly two lineages coalesce into one and that all lineages in the tree that are a distance \( d \) apart coalesce at the same time (perfectly balanced tree, figure 6).

An important difference between the corridor and the radial case is that in the radial case the MRCA of the whole population is always clearly identifiable. The symmetry relation in lineage densities highlighted in figure 5 then allows to build our tree forward in time from the centre of the colony as a branching process (rather than backwards as we did for the corridor).
Figure 6. Sketch of the deterministic tree model that captures that average lineage properties during radial growth.

We start with an arbitrary number \( b \) of starting branches, growing star-like towards the outside, splitting at deterministic distances away from the center based on the distance for which the branches have already been running, according to equation (7). Since we are in a radial setting with an inflation term, the distances between the lineages is dictated by both the angle between them, as well as the current distance away from the center.

After each splitting step, the angles are halved, leading to

\[
\alpha_i = \frac{\alpha_{\text{start}}}{2^i} \quad \text{for } i \in \{0, 1, 2, \ldots \}, \text{ with } \alpha_{\text{start}} = \frac{2\pi}{b}. \quad (8)
\]

For the very first step \((i = 0)\), where \(b\) branches start from the center, the height to splitting \(h_0\) is calculated as

\[
h_0 \sim (\alpha_0 h_0)^{3/2}
\]

\[
\Rightarrow h_0^3 \sim \alpha_0^3 h_0^3
\]

\[
\Rightarrow h_0 \sim \frac{1}{\alpha_0^3},
\]

where we used relation (7) together with the fact that in the very first step, the distance until the next splitting point is the same as the distance away from the center used for calculating the arc length distance between the lineages. For all following steps, the splitting heights are given via the equation

\[
h_i \sim \left( \alpha_i \left( \sum_{j=0}^{i-1} h_j + h_i \right) \right)^{3/2} \quad \text{for } i \in \{1, 2, 3, \ldots \}. \quad (10)
\]

For all \(i > 1\), equation (10) can be solved numerically. One finds that the equation always gives two positive solutions. As \(i\) increases, the next \(h_i\) is always dependant on the result of the previous heights \(h_0, \ldots, h_{i-1}\), requiring some considerations on the choice of solution. We find that choosing the bigger solution leads to increasing values of \(h_i\) until they diverge. Reciprocally, choosing the smaller solution leads to decreasing \(h_i\) that converge towards 0. It is important to point out here that, while we use the scaling in equation (7) to determine the location of the branching events, the relationship is expected to hold only when the distance between lineages is much smaller than the radius of the branching event \((d \ll r)\). For the first few steps this assumption is likely to break. However, since lineages double at every step while the time between steps grows at best as a power-law, the condition becomes true relatively quickly in the expansion process \((i > 2)\).

The simulation results for lineage density suggest that the distance between branching events becomes increasingly longer in the first half of the colony growth (up to \(R/2\)) after which it then becomes shorter and shorter with each subsequent step. As a result, for our model, we will choose the larger solution for the first \(p\) steps, and subsequently always choose the smaller solution. In the abstract model, the branching process can continue indefinitely, creating ever-shorter branches spaced closer and closer together as the tree grows outward. However, the successive values of \(h_i\) when \(i > p\) decline so rapidly that the sum \(\sum_{i=0}^{\infty} h_i\) converges to a finite value which corresponds to the radius \(R\) of the colony. Upon assembling deterministic trees with different values of \(p\), we observe that \(\sum_{i=0}^{p} h_i \approx \sum_{i=p+1}^{\infty} h_i \approx R/2\) independently of the value of \(p\). The geometry dictated by the solutions to equation (10) ensures that the transition from choosing the larger to the smaller solution always happens at approximately half of the colony radius, which coincides with the lineage density minimum in the Eden simulations (figure 5).

The universal tree model provides a deterministic prediction for the lineage density. As figure 6 shows, the number of lineages doubles at discrete values of the distance from the central node. By recording these values and the corresponding number of branches, we obtain a rescaled lineage density (symbols in figure 5) which reproduces the curves measured from Eden simulations. The behavior as \(x = r/R \rightarrow 1\) is dictated by the power law \(l(x) \sim s^{-2/3}\) expected from the corridor geometry and confirmed in figure 4: upon using the relation \(x = R(1 - x)\), we have \((r)/(2\pi r) \sim [R(1 - x)]^{-2/3}\). The universal tree recovers this scaling at large \(x\), but also reveals the behavior of the lineage density for \(r \lesssim R/2\), where inflation and branching play opposite roles (inflation creates space for lineages and branching events quickly fill it up). We find that the number of branches in the universal tree grows as \(l(r) \sim r^{1/3}\), leading to a divergence \(l(r)/(2\pi r) \sim r^{-2/3}\).
iors are well-captured by a phenomenological master curve

\[ \ln(\frac{l(r)}{2\pi}) = \frac{1}{\epsilon^{4/3} R^{2/3}} [x(1-x)]^{-2/3}, \]  

(11)

where \( \epsilon \) is a small-distance cutoff equal to the lattice spacing in the Eden model simulations (see appendix B for details). Equation (11) (dashed line in figure 5) successfully reproduces the lineage densities measured in Eden model simulations as well as the deterministic geometry generated by the universal tree model.

4.2. MRCA position of contiguous subsamples

The proposed master curve for the lineage density, equation (11), combines the effects of inflation and stochastic coalescence in a succinct form which explicitly captures the distinct tree structures for center distances below and above the value \( r = R/2 \). As an application of our results to a quantity of relevance to typical biological measurements, we now use this master curve to determine a general relationship for the radial distance \( R_{\text{MRCA}} \) of a contiguous sample of size \( n \) taken at the outer boundary of a colony that has reached a final radius \( R \).

We measure distances and sample sizes in units of the lattice spacing, and correspondingly set \( \epsilon = 1 \) in equation (11). If we assume a uniform angular distribution of lineages, then the number of surviving lineages for a contiguous sample of size \( n \) (covering an angle \( n/R \)) varies with distance \( r \) from the colony centre as

\[ l(n, r) = l(r) \frac{n}{2\pi R} \]  

(12)

Then the MRCA corresponds to the radius at which we are left with only one lineage \( l(n, R_{\text{MRCA}}) = 1 \), leading to the following equation

\[ x^{1/3}(1-x)^{-2/3} = R^{2/3} n^{-1}, \]  

(13)

which can be solved exactly. Because \( 0 < x = R_{\text{MRCA}}/R < 1 \), the acceptable solution to this equation is always unique and corresponds to

\[ R_{\text{MRCA}} = \frac{1 + 2R^2 n^{-3} - \sqrt{1 + 4R^2 n^{-3}}}{2R^2 n^{-3}}. \]  

(14)

For large \( n \), \( R_{\text{MRCA}} < R/2 \) and the scaling with sample size is such that \( R_{\text{MRCA}} \sim n^{-3} \). Conversely, if \( n \) is small, the MRCA is close to the edge on average, and its position follows the scaling \( 1 - R_{\text{MRCA}}/R \sim n^{1/3} \), which is analogous to the corridor case.

The transition between the two regimes corresponds to when \( R_{\text{MRCA}} \approx R/2 \). Using the equation above, we find that this corresponds to a critical angle (and critical sample size) \( \alpha_c = n_c/R \approx R^{-1/3} \). This scaling has been previously identified as the threshold frequency between bubbles and sectors in neutral mutations in two dimensional colonies [4].

4.3. Number of segregating sites

Analogously to the corridor case, the number of segregating sites \( S \) in a sample of size \( n \) is proportional to the total tree size that leads to the \( n \) surviving leaves, so that

\[ T_{\text{tot}}(n) = \int_{R_{\text{MRCA}}}^{R} l(n, r) dr \]  

(15)

\[ \propto nR^{1/3} \int_{R_{\text{MRCA}}/R}^{1} x^{1/3}(1-x)^{-2/3} dx \]  

(16)

\[ = nR^{1/3} \left[ B \left( \frac{4}{3}, \frac{1}{3} \right) - B \left( x_{\text{MRCA}}(n), \frac{4}{3}, \frac{1}{3} \right) \right] \]  

(17)

where \( l(n, r) \) are the number of lineages as a function of the distance from the center \( r \) that lead to the sample and \( B \) indicates the corresponding beta function.

4.4. Site frequency spectrum

Our tree model assumes a perfectly balanced tree with \( b \) initial branches that set the largest possible frequency of a mutation in the front population (1/b). Similarly to the argument for the corridor case, a mutation that is carried by at least \( n \) individuals at the edge has to occur somewhere in the tree before \( R_{\text{MRCA}}(n) \). Then, the number of mutations \( m(n) \) carried by at least \( n \) individuals at the edge is proportional to

\[ T'(n) \propto \int_{0}^{R_{\text{MRCA}}(n)} l(r) dr \]  

(18)

\[ \propto R^{1/3} B \left( x_{\text{MRCA}}(n), \frac{4}{3}, \frac{1}{3} \right) \]  

(19)

Figure 7 shows that the theoretical expectation of the cumulative site frequency spectrum \( 1 - \Pi(x) \), i.e., probability that at a mutation is carried by at least a fraction \( x = n/(2\pi R) \) of the population, without any additional fitting parameter, agrees remarkably well with the simulation. We observe a slight deviation at the point of inflection due to the discretized nature of the lattice in the simulations (the edge is only approximately one site thick). Interestingly, the agreement between theory and simulations suggests that in the radial case the tree is much more balanced then in the corridor case. The power-law tail, corresponding to an exponent of –4, is consistent with the site frequency spectrum of the full colony as we expect the large frequency mutations at the periphery to be stemming from sectors. The low frequency component of the site frequency spectrum is, in contrast, almost flat reflecting the fact that the later portion of the tree contributes negligibly to the total tree size.

In practice, often, only a subsample (or subsamples) of the colony periphery may be sequenced, as for instance in tumour biopsies [45, 54, 55]. Because of the spatial correlation of the genealogies, these subsamples can exhibit unusual signatures in the site frequency spectrum. We have found above that if two individuals (or samples) are picked farther than \( \alpha_c \) apart, their MRCA will very quickly converge to the centre of the colony. This implies the presence of long independent lineages that lead to the different samples over which several mutations can accumulate. Because these lineages do not branch for a long time, they will lead to a large number of mutations carried by a very specific frequency in the sample, showing
particular, 3D Eden modelsimulations have been shown to also of random-walk models outside the KPZ universality class. In easily applied to three dimensional growth and to other types.

evolutionary inference particularly challenging [45]. In this context, our analysis provides guidelines to design sampling schemes that can
display bubbles and sectors [4] analogously to the 2D case. We thus expect that a non-monotonic lineage density profile qualitatively similar to the one showed in figure 5, but characterised by different exponents, could be found also in this case, which is relevant to tumor growth.

Our results show, in agreement with previous work [43], that the lineages generated by an Eden model behave like directed polymers and can thus be modeled as random curves with super-diffusive statistics (mean-square transverse displacements grow faster than linearly with lineage length). The coalescence process (backwards in time) is then dictated by the annihilation of pairs of lineages as they collide. This analogy allows us to find a mathematical formulation for the average number of lineages as a function of time that lead to a final population at the edge of the expansion, which then can be used to provide estimates for the time to the MRCA and the number of segregating sites. Estimates for the site frequency spectrum require knowledge of the tree topology. Interestingly, here we find that the radial expansion is consistent with a balanced topology. By contrast, a linear front generates a site frequency spectrum inconsistent with a balanced topology, suggesting that rare long branches which coalesce well past the typical coalescence time play a crucial role.

While in this work we use the Eden model to describe the growth dynamics of two-dimensional populations building on previous studies [4, 20], this is by no means the only possible choice. The Eden model and the underlying KPZ universality class are relevant when the expanding front increases its roughness as it advances, due to a geometric feedback between local deviations from smoothness and the addition of new material at the front [38]. In some populations, however, front roughness might be suppressed and growth models that maintain a locally flat front are more appropriate. Alternative lattice models which maintain flat fronts, such as the Domany–Kinzel model [35] or flat-front stepping stone model [56] have been used to study clone dynamics in range expansions where front roughening is absent. The main difference compared with KPZ-type models is that the transverse wandering of lineages follows diffusive statistics (wandering exponent of 1/2) [43], and can be described by a Langevin equation with multiplicative noise [33, 35]. Our approach, with appropriately modified exponents, could be used to construct simplified genealogies for flat-front models of colony growth.

Recent studies have used both deep sequencing and lineage tracing techniques to generate a vast amount of data to disentangle tumour growth dynamics and selection [17, 19, 45, 54, 55]. In many cases, sector-like patterns are clearly observable suggesting that spatial growth and competition at the edge of the expansion play a crucial role. Because the cellular population is not well-mixed, results obtained from local sampling need to be carefully interpreted to infer the dynamics at play. For instance, the accumulation of mutations at specific frequencies, which in a well-mixed scenario would be interpreted a signature of selection, can be a sole consequence of the geometry of the sampling scheme, making evolutionary inference particularly challenging [45].
test whether a neutral spatial model is sufficient to reproduce the observed site frequency spectrum. Similarly, quantification of clone density as a function of time [17] and spatial location [19, 54, 55], and number of segregating sites [57] can provide orthogonal measurements to reveal whether selection or cell mobility is at play. Recent work has shown that these quantities can be more informative than the clone size distribution to identify selection in boundary-growing tumors [57], and our work provides analytical predictions for the neutral expectation, which can be compared with the experimental data.

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Data availability statement

All data that support the findings of this study are included within the article (and any supplementary files).

Appendix A. Computational methods

The computational data was obtained via simulations based on the Eden model [51]. The growth process in our simulations worked as follows: starting from a beginning set of cells, placed at fixed points on a lattice, we randomly select one cell, which still has at least one free, neighbouring lattice site. We then randomly choose one of these free neighbouring sites and place a new cell there. This new cell can be seen as the descendant of the initially chosen cell. We then repeat the process of randomly choosing a cell and placing a new cell at a neighbouring site, until we reach the desired colony size. An example of how an Eden growth process might look like, can be seen in figure 8. By tracking the information about each cell’s parent cell, we are able to trace back each cell’s lineage and obtain the location of the MRCA of two specific cells. The process for this is shown in figure 9.

Figure 8 shows how radial growth was initialized, where we started our simulation from a single cell. In contrast to this, the periodic corridor simulations start with a certain number of cells lying next to each other in a straight line. We then restrict the space that the cells can inhabit to a corridor of a width equal to the number of cells we start with, after which we start the Eden growth process as described above. The corridor boundaries are implemented to be periodical. In other words, the corridor could also be thought of as the outside surface of a cylinder, wrapping around and connecting to itself again. Therefore, a cell which is located at the edge of the corridor can give birth to a new cell at a free spot on the exact opposite side of the corridor. A visualization of this, as well as the growth process in a corridor in general, can be seen in figure 10.

In order to successfully sample MRCA data from random pairs of cells at the end of the corridor, all of these cells should be descendant of one single ancestor. In contrast to the radial colony, this is not automatically the cause though, since we start our corridor simulation with a column of unrelated cells (whereas in the radial case we start with only a single cell). Due to this, we have to let the colony grow for a significant amount of time until one lineage has pushed out all the others and the cells at the end of the colony are all descendant of one single cell from the very beginning of the corridor. The grid size required to handle corridors with sufficient length for this to happen greatly exceeds the computational memory limitations, even for small corridor widths. Due to this, we implemented a special method for simulating the corridor growth: in order to keep memory requirements low, we should only keep cells in memory which are part of the lineage of the cells at the very front, while deleting all other cells from the memory. At the same time we need to keep track of the spatial structure between the cells at the front, which is necessary for letting new cells grow using the Eden algorithm. To achieve this, we start with a small grid of pointers of the same width as our desired corridor, but a fixed length. Each of the starting cells in our first column has a distinct marker, which all of its descendants will also carry. When growing the grid, we do not save the cells’ information on the grid, but rather on an arbitrary data structure (e.g. a list) without any special spatial structure to it. The spatial structure for growing our colony is obtained through the small grid which points to the places in the memory where our cells’ information is kept (see figure 11). In figure 12 we sketched each step of the process for a better understanding. We start with a full column of cells, each with its own marker (see figure 12.1). The colony then grows until the first cells hits the right wall, at which point we stop the growth process. We then locate the most right column which is completely filled (marked with a black frame in figure 12.2) and check what markers are still present in that column. All cells with markers not present anymore can then be deleted from the data structure on the right (these cells have been circled in figure 12.2). After the unnecessary cells have been deleted, we can shift the corridor back to only include columns including and to the right of the column we just investigated for the remaining markers (see figure 12.3). From here, the process is repeated again, growing, sampling markers, deleting, until the sampled column only has markers of one kind. At this point we can delete the remaining other cells from our data structure. The colony has now successfully grown to the point that its front only includes cells which are descendant of one single common ancestor (see figures 12.4–7). At this point, we can stop growing our colony and start with analyzing its MRCA information.
Figure 8. Sketch of how a grid sequentially gets populated via the Eden model. The black dots indicate all possible spots where the new individual could be placed in the next step. The white arrows always point from parent to child.

Figure 9. Sketch of how lineages can be traced back to the MRCA. The lineages are traced back for two different cell pairs, one at the top of the colony and one at the bottom, indicated by the red dots. The lineages of the two cells in each pair are traced back, indicated by the red colouring of the lineage lines, until they intersect at the MRCA. The two MRCAs are indicated by the large red squares.

Figure 10. Sketch of an Eden model growth process in a corridor. From the upper left to the upper right picture the corridor gets populated, initially starting from a straight line of individuals filling the whole width of the corridor. The lower picture indicates how a cell at one edge of the corridor can have a descendant at the opposite side, due to the periodic boundary conditions of the corridor.

Figure 11. Sketch of the way that information is handled when growing large corridors. The square grid carries pointers to the actual data, which is kept in a data structure (marked in blue) without any spatial relationship between the individual data cells. The grid keeps track of the information about the spatial relationship between the individuals, while the other data structure carries the actual information of the individuals, like the information about their ancestor and children.

Appendix B. Lineage density from the universal tree model

Our considerations of the stochastic coalescences with statistics dictated by the KPZ wandering exponent (section 3), together with the incorporation of the inflation effect in the universal tree model (section 4.1), motivated the following scaling form for the lineage density of radial Eden clusters:

\[
\frac{l(r)}{2\pi r} = C[x(1-x)]^{-2/3}, \tag{B1}
\]

where \(x = r/R\) and \(C\) is an as-yet undetermined prefactor. The power-law divergences as \(x \to 0\) and \(x \to 1\) in equation (B1) are imposed by the KPZ scaling, but other considerations are needed to fix the value of \(C\).

The value of the prefactor \(C\) is determined by noting that the above scaling form is not valid out to \(x = 1\), but only up to some cutoff distance away from the outer limit of the colony whose value is set by the microscopic details of the growth process. Although the branching process of the universal tree model can be carried out to infinitely many steps, generating ever-finer leaves which approach the outer limit, the true coalescence process is limited by two microscopic length scales in any real biological system or realistic simulation thereof. First, genetic differences do not persist down to infinite resolution but instead are restricted to some finite spacing. In a microbial colony, for example, the smallest possible spacing...
between distinct genetic samples is the size of an individual cell; the lineages that are sampled are typically spaced even farther apart. Second, the KPZ wandering statistics arises within a coarse-grained description of the interface between the colony and its environment, which is only valid for roughness features that are larger than some microscopic length scale. This roughness scale is also of the order of a few cells for a microbial expansion.

To fix the value of the prefactor $C$, we impose these length cutoffs at the outer boundary of the colony. We denote the smallest spacing between distinct lineages by the variable $\delta$, and the smallest scale of roughness features by the quantity $\epsilon$. Therefore, the number of distinct samples at the outer boundary is $2\pi R/\delta$. Our proposed scaling form, equation (B1), is only valid out to distances within $\epsilon$ of the outer boundary and we do not expect any more mergers of lineages to occur between $r = R - \epsilon$ and $r = R$. To match the proposed master curve to the number of lineages at the outer boundary, we need

$$l(R - \epsilon) = C \left[ \left( 1 - \frac{\epsilon}{R} \right)^{-2/3} \right] \times 2\pi(R - \epsilon) = \frac{2\pi R}{\delta}.$$  

(B2)

Our scaling arguments are only valid provided $\epsilon/R \ll 1$. Keeping only terms to leading order in $\epsilon/R$ in the above equation, we find

$$2\pi RC \left( \frac{\epsilon}{R} \right)^{-2/3} = \frac{2\pi R}{\delta} \quad \Rightarrow \quad C = \frac{\epsilon^{2/3}}{R^{2/3}\delta}.$$  

(B4)

Equation (B4) specifies the prefactor by requiring the proposed master curve to match the lineage density at the edge of the colony. In the Eden simulations used in our work, both length scales $\epsilon$ and $\delta$ are given by the lattice spacing, hence we can set $\delta = \epsilon$ which gives

$$C = (R^2\epsilon)^{-1/3}$$

up to some $O(1)$ constant which we assume to be one. This value of the prefactor gives rise to the complete expression in equation (11). We find that the expression gives a reasonable match to the rescaled lineage density measurements from Eden simulations (compare dashed line to solid lines in figure 5) without any fits being performed. The agreement could be slightly improved by treating the $O(1)$ numerical constant as a free parameter whose value is determined by fitting the proposed master curve to the data. If we were to analyze genealogical tree data from a biological population, we would not have microscopic information about the quantities $\epsilon$ and $\delta$. In that...
case, it would be appropriate to use equation (B4) and fix the combination $e^{2/3}/δ$ as a fitting parameter.

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