Influence of Rough Flow over Sea Surface on Dry Atmospheric Deposition Velocities

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ABSTRACT

A Meteorological model and a dry deposition module were used to estimate the effects of sea surface rough flow (SSRF) over the sea surface on dry deposition velocities. The dry deposition turbulence resistance, $R_a$, and sub-layer resistance, $R_b$, decreased more than 10% and 5% due to SSRF, respectively. For example, for HNO$_3$, the mean dry deposition velocities ($V_d$) were 0.51 cm s$^{-1}$ in January, 0.58 in April, 0.65 cm s$^{-1}$ in July and 0.79 cm s$^{-1}$ in October with only smooth flow over the sea surface. However, the SSRF increased the $V_d$ of HNO$_3$ by 5 - 20% in the east China seas. These results show that SSRF is an important factor in estimating surface roughness to further improve calculation of the dry deposition velocities over the ocean. Improvements in parameterization of sea roughness length will be a worthwhile effort in related future studies.

Key words: Sea surface rough flow, Dry deposition, Deposition resistance, Numerical model

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1. INTRODUCTION

Terrestrial airborne pollutants are, as a matter of course, will be transported into the maritime atmosphere in different locales and deposited upon open sea surfaces. The deposition process depends on the turbulence of a marine atmospheric boundary layer and the characteristics of a sea surface. Much research has indicated that the roughness length is one of the main parameters which affect deposition processes and is of great importance for estimating the atmospheric deposition flux of a sea surface. Hertel et al. (1995) reported that dry deposition velocities of particles into the North Sea were lower than that upon land indicating a difference between land and water surfaces. However, the dry deposition velocities over an inland water surface usually were applied to calculate deposition fluxes to coastal water and open sea surfaces in many studies (Wan et al. 2002; Fu 2006). In fact, the deposition process on an open sea surface is very different from an inland water surface as characterized by wave and roughness lengths.

Early theoretical studies concerning an atmospheric deposition model focused on natural inland water surfaces. Slinn and Slinn (1980) built a deposition model suitable for natural waters considering the growth effect of particles with air humidity. Williams (1982) considered wave breaking, water droplets and the growth effect of particles in damp areas in a dry deposition model. In recent years, atmospheric deposition models have been developed and applied to estimate deposition flux over global and regional ocean surfaces (Asman and Janssen 1987; Andersen and Hovmand 1995; Ambelas et al. 2002; Jurado et al. 2004). Nho-Kim et al. (2004) calculated the dry deposition velocity of particles as a function of micro-meteorological conditions near the sea surface. Wang (2006) extended Williams’ dry deposition model for application to a sea surface and took into consideration of the effects of wind speed, humidity, broken surface coverage and transfer coefficients. Qi et al. (2005) adopted a particle growth formula developed by Gerber (1985) to improve the above model and applied it to calculate the flux of trace metals in particles into the sea in Qingdao.

However, many atmospheric deposition models still have limitations in sufficiently considering sea surface parameters. In an open marine boundary layer, sea surface roughness is not only the function of wind speed, or friction velocity, but also wave properties which can be calculated...
as a function of wave height, wave slope and wave age (Hsu 1974; Johnson et al. 1998; Smith et al. 1992). Taylor and Yelland (2001) considered the impact of waves on dynamic sea-surface roughness and built a parameterization scheme. Uz et al. (2002) conducted a series of flume experiments to study the wind-wave coupled system and calculated the contribution of waves to sea roughness length. Pan et al. (2005) confirmed the validity of Taylor’s scheme in calculating sea roughness using buoy data in marine sites and pointed out that the effects of wave properties on sea roughness should be taken into consideration under high wind speed. Nowadays, most studies on atmospheric deposition over water surfaces have mainly concentrated on particle, very few involved a gaseous substance. However, the contribution of gaseous substances to deposition fluxes to sea surface has been shown to be very important (Rendell et al. 1993; De Leeuw et al. 2003; Zhang et al. 2010). The gaseous dry deposition process is greatly influenced by meteorological parameters and its temporal and spatial variation is significant (Gao and Wesely 1995; Ma and Daggupaty 2000; Zhang et al. 2004, 2007).

In view of the limitations mentioned above, this study tries to explore the effect of sea surface rough flow (SSRF) on dry deposition velocities of gaseous substances using a meteorological model and a dry deposition module. The East China Sea has been taken as an example.

2. METHODOLOGY

2.1 Dry Deposition Model

During the dry deposition process, the constant flux layer can be divided into turbulent, viscous and surface layers according to deposition height (Walcek 1986; Wesely 1989). Gaseous dry deposition process is greatly influenced by meteorological parameters and its temporal and spatial variation is significant (Gao and Wesely 1995; Ma and Daggupaty 2000; Zhang et al. 2004, 2007).

For gaseous substances, Walcek et al. (1986) method is applied to calculate $R_d$; the land surface $R_s$ is calculated with the method used by Wesely (1989); and, the method developed by Hicks and Liss (1976) is adopted when calculating water surface $R_w$ as shown below.

$$R_s = \frac{\ln(Z_e / Z_0) \cdot \psi_c}{kU_c}$$  \hspace{1cm} (1)

$$R_b = \frac{2}{kU_c} S_c^{2/3} \quad \text{(land surface, Wesely 1989)}$$

$$R_w = \frac{\ln(kU_cZ_0/D_c)}{kU_c} \quad \text{(water surface, Hicks and Liss 1976)}$$  \hspace{1cm} (2)

where $Z_e$ is the height at which the dry deposition velocity is calculated, $Z_0$ stands for surface roughness length and is usually derived from empirical values according to localized topographic characteristics, $\psi_c$ is the revised stability function of mass. $U_c$ is the friction velocity; $S_c$ is the Schmidt number; and, $D_c$ is the molecular diffusivity.

$R_c$ is closely related to the biochemical reaction between pollutants and the deposition surface, which is dependent on the solubility, reactivity and the surface characteristics of the gas.

The dry deposition velocity $V_d$ can be calculated as:

$$V_d = \frac{1}{R_s + R_b + R_c}$$  \hspace{1cm} (3)

2.2 Parameterization of Sea Surface Roughness Length and Sensitivity Schemes

Sea roughness, also called sea surface aerodynamic roughness length, is defined as the height at which wind speed equals zero above the sea surface. It depicts the micro-scale sea surface roughness, and the change illustrates the main characteristics of momentum transfer between the atmosphere and ocean to some extent. Sea roughness can seldom be directly observed. Since wind speed is easy to observe, sea surface roughness usually is parameterized as a function of wind speed.

Charnock (1955) proposed the classic water surface roughness scheme as $Z_o = Z_{ch}U_{*}^2/g$, where $U_*$ is the friction velocity; $g$ is the gravitational acceleration; and, $Z_{ch}$ stands for the Charnock parameter which characterizes the water surface properties. The value range of $Z_{ch}$ differs for lakes, limited surf zones and broad water surfaces. For the calculation of deposition resistances, the water roughness scheme recommended by Hicks and Liss (1976) was based on Charnock relation. In the Meteorology-Chemistry Interface Processor (MCIPv3.2) for Models3 Modeling System (Byun and Ching 1999), the calculation scheme for roughness is also based on Charnock’s relation. In combination with Wu’s (1982) research, the parameterization of $Z_o$ is calculated as $Z_o = 0.0185U_{*}^2/g + 0.0001$ (m), where the first term stands for rough flow and the second term represents smooth flow.

Hence, this study has introduced the Taylor and Yelland (2001) scheme into the deposition module to calculate the rough flow on sea surface aerodynamics roughness which includes the effects of wave height:

$$Z_{0h} = 1200h_w (\frac{h_i}{T_p})^{4.5} + 0.0001$$  \hspace{1cm} (4)

$$h_i = 0.018U_{10}^2(1 + 0.015U_{10}), l_p = t_v c_w,$$

$$t_v = 0.729U_{10}, c_w = 9.8 t_v / 2 \pi$$  \hspace{1cm} (5)

where $Z_{0h}$ is sea surface roughness (m), $h_i$ is effective wave
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In this study, sensitivity tests were conducted under two different schemes for comparison to study the effects of sea surface rough flow on dry deposition processes. The sea surface roughness of Scheme one is fixed to 0.01 cm, only considering smooth flow. Scheme two additionally included a dynamic parameterization term consider sea surface rough flow (SSRF), as listed in Eq. (4).

The sensitivity parameter in the deposition model shows that roughness has a direct impact on the turbulent and viscous layer resistance. Deposition resistance has corresponding changes in accord with seasonal and spatial variation of roughness which results in the changes of dry deposition velocities. In Eq. (7), \( \Delta V \) refers to the change ratio of the variables, \( V_{\text{surf}} \) is the variable with rough flow, \( V_{\text{non}} \) is the variable with only smooth flow (roughness = 0.01 cm), where V can be referred as aerodynamic resistance \( (R_a) \), viscous layer resistance \( (R_v) \) and the dry deposition velocity of various pollutants \( (V_d) \).

\[
\Delta V = \frac{V_{\text{surf}} - V_{\text{non}}}{V_{\text{non}}} \tag{7}
\]

2.3 Meteorological Parameters

Lin et al. (2004) reported that the meso-scale meteorological model MM5 performs well in simulating a sea surface wind field. In their results, the simulation deviation in 10-m wind speed and wind direction above sea surface in the East China Sea was under 20%. Thus in this study, the meteorological fields in January, April, July, October 2007 were calculated by MM5 to drive dry deposition module. The Lambert projection was used in this simulation. The physical processes schemes including Grell’s cumulus parameterization scheme (Grell et al. 1994), MRF boundary layer scheme (Hong and Pan 1996), vapor scheme (Dudhia 1996), five-layer (1, 2, 4, 8, 16 cm) soil model and flexible boundary conditions were applied. The initial and boundary conditions were provided by 1° × 1° reanalysis data with a time interval of 6h from the European Centre for Medium-Range Weather Forecasts (ECMWF). Two domains were nested within the module as shown in Fig. 1. Domain 1 covers the entire land region of China and the East China Sea, with the center located at N35°, E105°. There are 55 × 75 grids in this domain with a 81-km grid distance. Domain 2 covers eastern China including the sea, with 70 × 70 grids and a 27-km grid distance. The terrain-following coordinate \( (\sigma) \) is defined as the vertical coordinate with non-equidistant grids spreading in 24 layers; nearly half of the grids are distributed below 2 km in order to better illustrate the structure of the atmospheric boundary layer. There were 24 land-use
types from the USGS Satellite Dataset with a spatial resolution of 2' (about 3.7 km) applied in the module.

Observational data of monthly mean of wind speed from two coastal meteorological sites, Lvsi (N32.06°, E121.60°) and Shengsi (N29.60°, E121.81°) were used to validate the 10-m height wind speed simulated by MM5 (Fig. 2). Generally, the modeled wind agreed well with the observed ones. The correlation coefficients reached above 0.65.

3. RESULTS AND DISCUSSION

3.1 Spatial and Temporal Distribution of Sea Surface Roughness in the Eastern China Sea

The spatial and temporal sea surface roughness has been estimated using the rough-flow parameterization scheme with the meteorological field as input data.

The mean monthly wind speed of each typical month varies greatly which is consistent with the results in the North-west and North Pacific Ocean, studied by Liu and Sun (2000) and Lin and Chen (2002), respectively. Generally, the average wind speed in fall and winter is greater than that of summer and spring. Figure 3 shows the distribution of mean sea surface roughness in typical months. In January, roughness length varied between 0.01 and 0.017 cm in most simulation regions with higher values of 0.02 - 0.022 cm occurring in the coastal region along Fujian Province due to integrated effects of cold high-pressure system and strait terrains. There were more significant spatial gradient in roughness in April and July, with southern high values in July and northern high roughness in April, both ranging from 0.011 to 0.018 cm. In October, the sea surface roughness in dominant ocean areas was from 0.019 to 0.032 cm, 2 to 3 times of the roughness of smooth surface and gradually decreasing from the southeast to northwest. In Bohai Bay, the roughness value was reduced to 0.011 to 0.022 cm. In the southeast coastal regions, the mean sea surface roughness in October was highest in all typical months, which is in accordance with the seasonal variation of sea surface roughness using TOPEX altimeter wind data (Zhou and Guo 2005). This could be because of local tropical storms and synoptic convergence systems driving the enhancement of wind speed in the coastal region in October.

![Fig. 2. Comparisons of simulated and observed 10-m wind speeds (a) Lvsi site, (b) Shengsi site.](image-url)
3.2 Impact of Sea Surface Rough Flow on Dry Deposition Resistances

The aerodynamic resistance $R_a$ and viscous layer resistance $R_b$ in only smooth flow situation were shown in Table 1 as reference background to show the impact of rough flow on resistances. Impacts on aerodynamic resistance $R_a$ and viscous layer resistance $R_b$ were estimated when the variable $V$ in Eq. (7) is replaced by $R_a$ and $R_b$, respectively.

Statistically, 1745 grids in a simulation domain were taken as sea region samples to estimate change ratios for $R_a$ and $R_b$ in different sea surface areas after considering SSRF (as in Fig. 4, only January and July were presented here). For $R_a$, the monthly average decrease were primarily distributed between 0.08 and 0.11 and decreased by around 10% in more than 80% of the sea surface in January. However, the change ratio in April is less than 0.10 within most sea surface, only 20% sea surface was reduced by more than 0.10. The change ratio in July has an inhomogeneous spatial

Fig. 3. Distribution of mean sea roughness lengths in typical months (cm) (a) January, (b) April, (c) July, and (d) October.
distribution which varies from 0.05 to 0.16; but, each 2% change ratio section covered about 4% sea surface. The change situation in October distributed the widest range 0.05 - 0.19 showing a significant spatial distribution, and $R_a$ in more than 60% of the sea surface decreased by more than 14%. The seasonality of the $R_a$ change ratio was the result from the integrated impact of the change of sea surface roughness and friction velocity.

For $R_b$, the percentage of change ratio mainly ranged from 0.025 - 0.04 in January and in 0.01 - 0.04 in April. The highest change ratio in July and October reached 0.06 and 0.07, respectively and the change ratio was mainly concentrated between 0.05 and 0.07 in October. All the change ratios in typical months were below 0.07 for $R_b$ which endured less impact by rough flows compared to $R_a$. In addition, as shown in Table 1 that $R_a$ is 4 - 5 times of $R_b$. Thus, the change in $R_a$ would account for the dominant contribution to the change in other resistances.

### 3.3 Temporal and Spatial Distribution of Dry Deposition Velocity and Its Response to Sea Surface Rough Flow (SSRF)

The deposition velocity $V_d$ of HNO$_3$ is decided primarily by the aerodynamic ($R_a$) and viscous layer ($R_b$) resistance
since surface resistance ($R_c$) of HNO$_3$ is rather small and it could be ignored during its deposition process to water surfaces (Walcek 1986; Zhang et al. 2004). Thus, HNO$_3$ has been chosen as an example to explore the impact of SSRF on dry deposition velocity.

The dry deposition velocities of $V_d$ of HNO$_3$ using smooth-flow scheme were presented in Figs. 5a and b. The dry deposition velocities $V_d$ of HNO$_3$ distributed mainly between 0.5 to 4.5 cm s$^{-1}$ over land regions. The modeled $V_d$ over land is in accord with the observed $V_d$ of 0 - 4.7 cm s$^{-1}$ over grassland, 2.2 - 6.0 cm s$^{-1}$ and 0.8 to > 20 cm s$^{-1}$ for mixed forest (Huebert and Robert 1985; Meyers et al. 1989; Sievering et al. 2001). The deposition velocity of HNO$_3$ to water surface is smaller than that to land surface with stronger turbulent activity above plant canopies. The $V_d$ varied from 0.3 to 0.7 cm s$^{-1}$ on the eastern China sea surface, and higher in summer with the stronger turbulent activity making deposition easier than that in winter. The modeled $V_d$ over

![Fig. 5. $V_d$ of HNO$_3$ in (a) January and (b) July (cm s$^{-1}$) and change ratio in $V_d$ in (c) January and (d) July (%).](image-url)
smooth sea flow was comparable with 0.26 to 0.42 cm s\(^{-1}\) over water surface estimated by Matsuda et al. (2001).

\(V\) in Eq. (7) was replaced by \(V_d\), the change ratio of \(V_d\) was calculated after considering rough flow and presented in Figs. 5c and d. The change ratio depends on both the sea surface roughness and \(V_{\text{dmean}}\). Thus, the areas, where largest changes in \(V_d\) (Figs. 5c and d), did not wholly cover the areas with largest surface roughness (Fig. 3). However, \(V_d\) still has an obvious response to sea surface rough flows with space and season. The spatial distribution of the change ratio showed that \(V_d\) increased by 7 - 10% in most ocean areas and more than 13% in the middle East China Sea in January. \(V_d\) in July increased by 8 - 16% in most ocean regions. In October, \(V_d\) increased by more than 17% in most of the south-east parts of the East China Sea, even up to 20% in North-West sea regions to Taiwan island (figure was not presented here).

In addition to HNO\(_3\), the monthly mean change ratios of \(V\) for NO\(_2\), NO, and NH\(_3\) were listed in Table 2. Similar to HNO\(_3\), the impact on \(V\) of NH\(_3\) from rough flows was also significant, with the annual mean increase rate of 10.1%. However, NO\(_2\) and NO receive other resistances despite \(R_a\) and \(R_v\) during deposition to water surface, which reduces the impact of rough flow on \(R_a\) and \(R_v\), thus limits the impact on their dry deposition velocity.

### 4. CONCLUSIONS AND DISCUSSION

This study has introduced a sea surface roughness scheme with surf parameters into a dry deposition model to study the impact of rough flow on atmospheric dry deposition. The sea surface roughness increased from 0.01 to 0.03 cm after considering rough flow in the eastern China seas. Correspondingly, the mean value of aerodynamic \(R_a\) and viscous layer \(R_v\) resistances decreased 10% and 5%, respectively. The mean dry deposition velocities of HNO\(_3\) with only smooth flow in the eastern China sea were 0.51, 0.58, 0.65 and 0.79 cm s\(^{-1}\) in January, April, July and October, respectively. Correspondingly, \(V_d\) increased 9.4, 8.0, 9.5, 13.9% after considering SSRF. The impact of the SSRF on the dry deposition velocity has been proved to be of importance. Hence, sea surface dynamic roughness is an important parameter studying gaseous deposition above an ocean area and its parameterization scheme needs further improvement in atmospheric deposition model.

However, since the dry deposition process is also influenced by the characteristics of substances and other factors, impacts of sea rough flow on \(V_d\) vary with substance species. Substances such as NO\(_2\), NO are harder to deposit onto a sea surface and are less influenced by aerodynamic \(R_a\) and viscous layer \(R_v\) resistance, therefore receive smaller impact from SSRF.

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