Design Optimization of a Novel Networked Electromagnetic Soft Actuators System Based on Branch and Bound Algorithm

NAFISEH EBRAHIMI1, TEJA GUDA2, MILTIADIS ALAMANIOTIS3, (Member, IEEE), DIMITRIOS MISERLIS4, (Member, IEEE), AND AMIR JAFARI1, (Member, IEEE)

1Advanced Robotic Manipulators (ARM) Laboratory, Department of Mechanical Engineering, The University of Texas at San Antonio (UTSA), San Antonio, TX 78249, USA
2Department of Biomedical Engineering and Chemical Engineering, The University of Texas at San Antonio (UTSA), San Antonio, TX 78249, USA
3Department of Electrical and Computer Engineering, The University of Texas at San Antonio (UTSA), San Antonio, TX 78249, USA
4Department of Vascular Surgery, University of Texas Health at San Antonio, San Antonio, TX 78229, USA

Corresponding author: Amir Jafari (amir.jafari@utsa.edu)

This work was supported in part by the National Science Foundation (NSF) under Grant 1840834.

ABSTRACT This paper presents a design optimization framework based on Branch and Bound Algorithm for novel networked electromagnetic soft actuators. The soft actuators work based on the operating principle of solenoids but are made of intrinsically soft materials. We confirmed that by scaling down the size of the soft actuators, their force to volume ratio increases. This suggests that by miniaturizing the actuators size and attaching them as a network based on the arrangement of actin and myosin filaments in skeletal muscles, the output force can be enhanced. In order to achieve the maximum output force, design parameters of a single soft actuator as well as those of a network are considered as design variables. The maximum available volume (thickness, width and length) and deflection range of the network are considered design constraints. The cost function, i.e. the output force is a non-linear mixed-integer function. A Branch and Bound optimization algorithm based on interval analysis is then proposed to solve the optimization problem. Numerical simulations are presented for a representative example of an active soft brace for the human elbow joint. The results suggest that an electromagnetic soft actuator network can provide sufficient torque to be used as a drive train for an active elbow brace for both flexion and extension over a range of around 93 degrees.

INDEX TERMS Actuators, branch and bound, design optimization, elbow joint, electromagnetic.

I. INTRODUCTION

In the context of physical Human-Robot Interaction pHRI, the technology requirement to fabricate the robotic platform is fundamentally different than those developed for industrial robots [1]. The fact that pHRI robots, such as industrial exoskeletons, prosthetic devices and surgical robots physically interact with the soft tissues in the human body, exerting forces outside the tissues structural limitations, introduces medical hazards. Current industrial biomedical robots are designed for fast and accurate position control applications in a perfectly controlled environment without taking into account the force interaction uncertainties that exist in pHRI applications [2]. Consequently soft pHRI robots can provide a novel and safer option for biomedical applications [3], [4].

Actuators, as the source of producing force into robotic platforms, plays critical role in realizing the softness of the robot. As a result, soft actuators are gaining interests among the researchers and consequently many different types of soft actuators have been developed [5]. Soft pneumatic actuators [6]–[8] have been the first generation of soft actuators where the air pressure inside a soft and stretchable body, usually made of Polydimethylsiloxane (PDMS) [9], can be adjusted and inflate-deflate the robot’s body to create a desired type of motion and amount of force. This type of soft robots are fast in response, can generate considerably high amount of force and displacement. However, the air pressure has to be supplied by an air pump which is rigid and the
robot’s body has to be tethered to be attached to the pump, which not only makes the entire robot’s body “not soft” but also occupies large space if the source of power is taken into account.

Another type of soft actuators are based on Dielectric materials [10]–[13]. This type of elastomers are smart material systems that produce large strains. They belong to the group of electroactive polymers. Dielectric actuators transform electric energy into mechanical work. They are lightweight and have a high elastic energy density. However, they usually need extensively high amount of voltage (in order of 20KV) to be activated.

Shape memory actuators are another type of soft actuators that have been developed by researchers [14], [15]. A shape-memory alloy (SMA) is an alloy that can be deformed when cold but returns to its pre-deformed (“remembered”) shape when heated. Shape memory alloys are popular as actuators for soft bioinspired robots because they are naturally compliant, have high work density, and can be operated using miniaturized on-board electronics for power and control. However, SMA actuators typically exhibit limited bandwidth due to the long duration of time required for the alloy to cool down and return to its natural shape and compliance following electrical actuation.

Liquid-gas transiting actuators [16], [17] are another type of soft actuators, where a liquid with low boiling point (usually around room temperature) such as ethanol, encapsulated inside a soft and stretchable body such as PDMS. By increasing the temperature, through generating heat (electric current through a resistive tiny wire) the liquid would change its stage to gas and as a result huge changes in volume will happen. This would lead to inflation of the soft PDMS and thus motion and force can be generated. However, this type of actuators have very limited bandwidth as it takes time to reverse the motion by cooling down the gas into liquid.

Recently, some soft actuators have been developed based on the electromagnetic working principle of traditional rigid actuators [18]–[20]. In this type of actuators, the force is generated based on the Lorentz law: a conductive element carrying electrical current in a magnetic field will experience a force acting perpendicular to both the magnetic field and direction of electric current. This type of soft actuators do not need high amount of input voltage as Dielectric actuators do, and are fast in response and can develop considerably high amount of deflection. However, the challenge in these type of actuators is the relatively low amount of generated force. The reason is due to several factors: first of all, these actuators use a conductive liquid (usually Eutectic Gallium Indium or EGaIn) which have higher electrical resistance as compared to copper wires in traditional rigid actuators. Also, these actuators usually use flexible permanent magnets made by a mixture of PDMS and magnetic particles, which have lower magnetic strength as compared to rigid permanent magnets in traditional rigid actuators. Another challenge is due to limitations on fabrication process of the micro-channels inside PDMS. Some 3D printers are able to create very small size micro-channels, however, these devices are very expensive. A Comparative analysis of fabrication methods for achieving micro-channels in PDMS is presented by [21].

We have developed a novel Electromagnetic Soft Actuator (ESA) based on the working principle of solenoids with permanent magnets [22]–[24]. In our ESA, electromagnetic field is created by applying electric current through a soft conductive coil. The coil is made of PDMS micro-pipe with diameter of around 0.1mm filled by EGaIn. In our design, two coils are antagonistically embedded inside a PDMS body with a springy connection in between as it is shown in Figure (1). Once electric current is being supplied the two coil can get magnetized based on the Lorentz law and attract each other. In order to intensify the electromagnetic field between the two coils, a flexible permanent magnet is placed inside the coils.

The permanent magnet is made by pouring a mixture of PDMS and magnetic particles into a 3D printed cylindrical-shaped mold, where the mixture in placed inside a strong external magnetic field during the curing process, to align the magnetic orientation of each particle. As a result, once the mixture is completely cured, the magnetic particles remain aligned even when the external magnetic field is removed. The result is a cylindrical flexible material with magnetic properties.

This type of soft actuator is highly scalable where the scaling factor is determined by the available manufacturing technology. For example, the smallest PDMS micro-pipe available has a diameter around 0.1mm, whereas using high-end 3D printers, PDMS can be printed as a coil-shape with embedded helical micro-channel inside with diameter in the order of micrometers. However, due to extremely high price of such 3D printers, this manufacturing technology is very expensive and only justifiable in mass production.

![Design concept of an Electromagnetic Soft Actuator ESA: coil is composed of micro-channels by wrapping flexible wire around the core. Once electric current flows inside the channels it creates electromagnetic field. Two opposite magnetic poles attract each other and create force that deflects the spring linkage between the two coils.](image-url)
By decreasing the size of the actuator, the generated force will be reduced. Interestingly we found that by scaling down the ESA’s size, its force to volume ratio increases. This was confirmed through experiments on different sizes of ESAs [24]. This behaviour was also analytically proven using the geometry of the actuator and electromagnetic equations based on the Lorentz Law [25].

This property of ESAs suggests that by reducing the size of ESAs and attaching them in series and parallel as a network, the output force can be enhanced compared to a single ESA with the same size of the networked ESAs. Due to their light weight, their fast response, and being able to be operated with voltage range around 10V 80V, with increased output force, ESAs can reproduce a soft actuation technology suitable for pHRI, especially for rehabilitation and support of patients with mobility impairments.

As it will be shown in section 3, the equation governing the force is nonlinear and contains design variables that are both real and categorical in nature. In order to maximize the output force, therefore, we are dealing with a nonlinear mixed-integer optimization problem. We propose to use a Branch and Bound algorithm based on interval analysis technique that is capable of obtaining the global optimal with mixed-integer variables.

In order to examine the capability of a networked ESAs to be used as drive train for a rehabilitation or force augmentation device, we consider a case of an active elbow brace and use of our proposed Branch and Bound optimization framework to determine whether or not our optimal network of ESAs can achieve the required amount of force and deflection.

The organization of this paper is as follow: in section 2, the design parameters will be explained and the output force of a single actuator as well as that of a network will be formulated as a function of design parameters. In section 3, the proposed Branch and Bound design optimization algorithm based on interval analysis will be presented. Section 4, presents an active elbow brace as a case study and applies the proposed optimization algorithm to determine the capability of networked ESAs as a drive train for such a device. Finally section 5 presents the conclusion and future works.

II. DESIGN PARAMETERS OF NETWORKED ELECTROMAGNETIC SOFT ACTUATORS AND FORCE FORMULATION

In this section, we first formulate the force as a function of design parameters of a single and then double (antagonistic) coils. We then include the effect of magnetic core to calculate the resultant force. In the next step, the generated force of networked ESAs will be formulated as a function of its design parameters.

A. ELECTROMAGNETIC FORCE OF TWO ANTAGONISTIC COILS

As mentioned before, our proposed ESA consists of two electromagnetically inductive coils that contain a flexible permanent magnet core. The force generated by the coils depends on several parameters as following: number of turns, number of loops, length and diameter of the coils, electrical current, length and diameter of the flexible permanent magnet, and the type and ratio of mixture of materials being used.

To design the ESA, a theoretical analysis should be conducted on the electromagnetic force that is being produced by the two coils, taking into account that the two coils share a common flexible permanent magnetic core. For this purpose, we first calculate the magnetic field by the Biot–Savart Law [26] along the axis of a representative loop of the coil with unit vector of $\hat{k}$, considering an arbitrary point $P$ through which a steady current $I$ is passing (Figure (2-a)). The magnetic field at point $P$ due to the contribution of the current element can be presented by applying Biot–Savart Law, as follow:

$$\mathbf{d}\mathbf{B} = \frac{\mu_0 I dS}{4\pi} \times \mathbf{r}_e.$$ (1)

where $\mu_0$ is the permeability constant synonymous to the permeability of free space or as magnetic constant, $\mathbf{r}_e$ is the effective radius defined as:

$$\mathbf{r}_e = \frac{R}{\sqrt{\rho^2 - \rho'^2}}, dS = R d\theta \hat{r}, \mathbf{r}_e' = \hat{r}.$$ (2)

$\mathbf{r}' = R\hat{r}$ and $R$ is the magnitude of the coil radius vector $\mathbf{r}'$ as it is shown is Figure (2-b).

![FIGURE 2. a): Magnetic field at the desired point P resulted by a single circular current element and b): Magnetic field at a desired point generated by a schematic coil.](image)

To calculate the magnetic field at point $P$, $d\mathbf{B}$ has to be integrated over the entire circular loop. The first integral vanishes due to the fact that radial unit vectors around the circle sum to zero. Hence, the overall would only be the axial component of the magnetic field as:

$$\mathbf{B}_z = \frac{\mu_0 R^2}{4\pi (Z^2 + r^2)^{3/2}} \int_0^{2\pi} d\theta \hat{k}.$$ (2)

This is the magnetic field of a single loop acting on point $P$ along the axis $\hat{k}$. A solenoid consists of many loops and layers of loops. In order to calculate the magnetic field acting on an arbitrary point $P$ along the axis $\hat{k}$ in a solenoid, the resultant magnetic field of a single loop should be integrated over the entire length of the Solenoid. $dz'$, as shown in Figure 2, is a selected current element on the solenoid. The $z$ dimension of the selected point $P$ is always measured from the central loop of the solenoid, where the amount of current $dl$ passing
through the element \(dc'\) is \(nldc'\), where \(n = \frac{N}{A}\) is the coil’s loop density.

The infinitesimal magnetic field due to the presence of an infinitesimal electric current within the length of \(dc'\) is:

\[
\overrightarrow{dB} = \frac{\mu_0 R^2}{2[(z - z')^2 + r^2]^{3/2}}dl. \tag{3}
\]

Similarly by integrating \(\overrightarrow{dB}\) over the entire length of the solenoid, the total magnetic field at point \(P\) can be found as:

\[
\overrightarrow{B} = \frac{\mu_0 n I}{2} \left[ \frac{1/2 - z}{[(z - l/2)^2 + R^2]^{3/2}} + \frac{1/2 + z}{[(z + l/2)^2 + R^2]^{3/2}} \right]. \tag{4}
\]

This is the magnetic field of a solenoid at any given point \(P\) along its central axis \(k\) with a distance \(z\) from its center that is located along the axis \(k\) at the middle of solenoid’s length.

In our ESA’s design, two antagonistic coils are interacting with each other and therefore, contribution of both coils on a given point should be taken into account. By calculating magnetic field due to each coil (left and right) and then applying the superposition principle, we can find the total magnetic field at any given point \(P\) located on the common central axis of the two antagonistic coils.

\[
\overrightarrow{B}_{\text{ext}} = \overrightarrow{B}_{\text{right}} + \overrightarrow{B}_{\text{left}}. \tag{5}
\]

To calculate each magnetic field, the distance of point \(P\) from the center of its corresponding coil \((z_1\) and \(z_2)\) should be calculated.

### B. EFFECT OF PERMANENT MAGNET CORE

In order to determine the effect of flexible permanent magnet core on the resultant force, we will establish the Electromagnetic Charge Model. The Charge Model is a useful method for analyzing the force of permanent magnet. Given \(\rho_m = -\nabla \cdot \overrightarrow{M} (A/m^2)\) as equivalent volume charge density and \(\sigma_m = \overrightarrow{M} \cdot \hat{n} A/m\) as equivalent surface charge density (\(\hat{n}\) is the unit vector of the surface), the resultant force due to presence of a permanent magnet can be calculated as:

\[
F = \int_v \rho_m \overrightarrow{B}_{\text{ext}} dv + \int_v \sigma_m \overrightarrow{B}_{\text{ext}} dv. \tag{6}
\]

The permanent magnet has a consistent and uniformly distributed magnetic field along its central axis as \(\overrightarrow{M} = M\hat{z}\) which makes \(\rho_m\) equal to zero. The net magnetic moment per unit volume of the permanent magnet is \(\overrightarrow{M}\) that can be calculated as: \(\overrightarrow{M} = B_s/\mu_0\), where \(B_s\) is remanence or residual flux density that is obtainable from the hysteresis curve of the material.

The unit vector of the surface \(\hat{n}\) has three distinct directions as follow: when \(z = 0\), \(\hat{n}\) is opposite of \(\hat{z}\), when \(z = h\) \((h\) is length of the cylinder), \(\hat{n}\) is equal to \(\hat{z}\) and when the distance from the central axis \(r\) is equal to radius of cylinder \(\overline{r}\), \(\hat{n}\) is perpendicular to \(\hat{z}\). Therefore, if the magnetization of the magnetic core along \(z\) axis is \(M_z\), the surface charge density \(\sigma_m\) would be \(-M_z\) when \(z = 0\) and \(M_z\) when \(z = h\).

However, of the surface of the permanent magnet where \(r = \overline{r}\), the surface charge density would be zero as the two vectors \(\hat{z}\) and \(\hat{r}\) are orthogonal.

In our ESA design, as it is shown in Figure 3, the North pole of the permanent magnet is located in the right coil while the South pole is in the left coil. The dimension \(d\) is the distance between each coil’s mid point and the magnetic poles. Each pole has the same distance to the correlated coil’s end as the magnetic core is placed precisely at the middle of the line that is connecting the two coils. Due to symmetrically antagonistic geometry of our design, only the applied force at one end surface of the magnetic core has to be calculated and then doubled to obtain the total applied force.

### FIGURE 3. Configuration of ESA with two schematic coils and a common magnetic core.

In order to find the magnetic field at any arbitrary point within this design, let’s consider \(z_1\) and \(z_2\) to be \(d\) and \(d + h\), respectively. The total magnetic field due to left and right side coils can then be calculated using Eqs. 4 and 5 as:

\[
\overrightarrow{B}_{\text{ext}} = \frac{\mu_0 n I}{2} [B_1 + B_2 + B_3 + B_4]. \tag{7}
\]

where

\[
B_1 = \frac{1/2 - d}{[(-d + l/2)^2 + R^2]^{0.5}}, \quad B_2 = \frac{1/2 + d}{[(d - l/2)^2 + R^2]^{0.5}},
\]

\[
B_3 = \frac{1/2 - (d + h)}{[(d + h) - l/2]^2 + R^2]^{0.5}}
\]

and

\[
B_4 = \frac{1/2 + (d + h)}{[(d + h) + l/2]^2 + R^2]^{0.5}}.
\]

The resultant force can be formulated using Eq. 6 as:

\[
F = \overrightarrow{B}_{\text{ext}} M \int_0^\overline{r} \int_0^{2\pi} r dr d\phi = \frac{\overrightarrow{B}_{\text{ext}} B_s}{\mu_0} \pi \overline{r}^2. \tag{8}
\]

By replacing Eq. 7 into Eq. 8, the force that applied to either North or South Pole of the magnetic core and generated by magnetic fields of both the right and the left coils can be calculated as:

\[
F = \frac{nIB_s}{2} \pi \overline{r}^2 [B_1 + B_2 + B_3 + B_4]. \tag{9}
\]

Therefore, the total force applied to the magnetic core both at the North and South poles would be \(F_{\text{total}} = 2F\). The current in this equation can be determined by the input voltage \(V_{in}\) and equivalent resistance of the coil’s circuit \(R_{eq}\) as \(I = \frac{V_{in}}{R_{eq}}\). The equivalent resistance of the coil is a function of length \(L_w\) and cross-section area \(A_{ws}\) of its conductive wire.
as well as specific resistivity of the conductive material $\rho_w$ as $R_{eqv} = \frac{\rho_w}{\pi \mu}$. The total force a nonlinear function in which we are dealing with both real variables such as dimensions as well as categorical variables such as the residual flux density or specific resistivity of the materials used in ESA.

Figure 4 shows components of a representative single ESA and its flexible magnetic core.

![FIGURE 4. a) Elements of ESA; hollow cylinder with embedded conductive coils, b) soft magnetic core and c) flexibility of ESA.](image)

C. NETWORK OF ELECTROMAGNETIC SOFT ACTUATORS

A network of ESAs consists of ESAs attached in series and parallel, similar to the arrangement of Actin and Myosin filaments in skeletal muscles. When ESAs are attached in series fashion, they compose a Fiber and when fibers are attached in parallel, they make a Fascicle. Output force of a fiber is equivalent to the output force of a single ESA while the total deflection of a fiber is summation of deflections of its ESAs.

The output force of a fascicle, however, is the summation of output force of its fibers, which, as mentioned before, is equivalent to the summation of output force of single ESAs attached in parallel. Therefore, to maximize the output force of a networked ESAs, in addition to maximizing the output force of a single ESA, the parallel arrangement of ESAs should be optimized. This parallel arrangement is a function of the cross-section profile of ESAs. Usually coils have circular cross-section profile, but the conductive wire can also be wrapped in different cross-section shapes such as triangular, square, etc. basically if $n$ represents number of sides of a shape, the cross-section profile can have $n$ as triangular, square, etc. basically if $n$ represents number of sides of a shape, the cross-section profile can have $n$ sides where $n$ can change from 3 (i.e. triangle) to $\infty$ (i.e. circle).

The purpose of this subsection is to geometrically design a fascicle network of ESAs so that the magnetic field and force out of a certain occupied volume of actuator network can be maximized. In the previous subsection we formulated the magnetic field and output force of a single ESA mainly by taking into account the geometry of the coils embedded in it. Now, we intend to discuss the formulation of a network of soft actuators, which are positioned with their long axis parallel next to each other. Given the point that the coil is the heart of the proposed ESA, we need to study the coil’s cross-section profiles. The idea is to eventually find the best cross-section profile for the axially packed network of solenoids shown in Figure 5 which yields to the maximum generated magnetic field and force out of a certain occupied volume of the network. Since, we intend to compare cross section profiles for each solenoid profile, we just need to calculate the magnetic field and force of one loop coil.

![FIGURE 5. Some examples of cross-sections of electromagnetic actuator networks with various coil profiles including: triangular, square, pentagonal, hexagonal, octagonal and, circular shapes.](image)

The polygons and circle shapes for the coil’s section profiles will be studied here. For this purpose, similar to the approach was taken in previous subsection, the magnetic field of the considered profile needs to be determined using Biot-Savart Law. For all the polygonal profiles, firstly, we need to calculate the magnetic field generated by a finite straight line representing the side of the polygons at a desired point as shown in Figure 6.

![FIGURE 6. Magnetic field of a finite straight wire at the desired point $P$.](image)

The current carrying element in $x$ direction is $d \vec{S} = dx \hat{i}$. The goal is to find the magnetic field of the current carrying wire at point $P$. The position vector describing $P$ is $\vec{P} = a \hat{j}$. The relative position vector which points from the source point to the field point is $\vec{r} = \vec{P} - \vec{r}$. For point $P$ with the location of $(0, a)$, the relative position vector would be $\vec{r} = a \hat{j} - x \hat{i}$. The magnitude of this vector defines the distance between the current element and the point $P$ and is equal to $r = (a^2 + x^2)^{0.5}$. The corresponding unit vector is given by $\hat{r} = \frac{\vec{r}}{r} = \sin \theta \hat{j} - \cos \theta \hat{i}$.

According to the Biot-Savart Law the contribution to the magnetic field due to the current carrying element is:

$$d \vec{B} = \frac{\mu_0 I d \vec{S} \times \hat{r}}{4\pi r^2}. \tag{10}$$

Substituting $a = r \sin \theta$ and $dx = -a \csc^2 \theta d\theta$ in Eq. 10, the differential contribution on the magnetic field is obtained as:

$$dB = \frac{\mu_0 I}{4\pi a} \sin \theta d\theta. \tag{11}$$
Integrating this infinitesimal magnetic field over the range of angles from $-\theta_1$ to $\theta_2$, the total magnetic field can be calculated as:

$$B = \frac{\mu_0 I}{4\pi a} \int_{-\theta_1}^{\theta_2} \sin \theta d\theta = \frac{\mu_0 I}{4\pi a} \cos \theta_2 - \cos \theta_1. \quad (12)$$

In case of symmetrical arrangement of the conductive wire around perpendicular line passing through point $P$ (i.e. $\theta_1 = \theta_2 = \theta$), the total magnetic field would be:

$$B = \frac{\mu_0 I}{2\pi a} \cos \theta. \quad (13)$$

Now that we calculated the magnetic field of a current carrying finite wire, it is possible to calculate the magnetic field produced by polygon shaped current caring loops.

For each polygonal shapes, the magnetic field at the center of the section is obtainable by superposition of the magnetic field produced by each sides at the center of the profile. For instance, for a triangular profile, depicted in Figure 7 the magnetic field at the center of the section would be:

$$B_{tri} = 3B_{line} = \frac{3\mu_0 I}{2\pi a_{tri}} \cos \theta_{tri}. \quad (14)$$

where $a_{tri}$ is the apothem of the triangle. Therefore, in order to compare various cross-section geometries, we would equalize different sections in terms of cross-section area and find the correlated edge length of each polygon solenoid section. In other words, we assume that all polygon sections have same cross-section area and then the side’s length for each one can be calculated. The cross-section area can be calculated by the apothem of each polygon and number of the sides on each polygon as: $A_{cross-section} = m \cdot a^2 \cdot \tan(180/\text{m}).$

Once, the magnetic field of all cross-section profiles has been calculated, the generated force resulted from the interaction between coil and magnetic core can be obtained:

$$F = B_{coil} \frac{B_r}{\mu_0} A_{cross-section}. \quad (15)$$

For each coil profile geometry, the magnetic core has the same profile with the correlated coil. Since, we assumed earlier that all profiles have the same cross section area, the magnetic core cross-section $A_{cross-section}$ would be the same as well. Therefore, the produced force from the coil would be proportional to its magnetic field for each case, hence the ratio of generated forces would be the same as the ratio of magnetic fields.

Another parameter to be considered is packing density $\lambda$ that can be defined as a function of total cross-section area of the network $A_{network}$ and cross section area of a single ESA and number of parallel fibers within a network $P$ as: $\lambda = PA_{cross-section}/A_{network}$. If all cross-section area of the network can be occupied by ESAs, then $\lambda = 1$.

Let’s say number of series ESAs within a fiber is $S$ and number of parallel fibers within a network is $P$. Then the total elongation of the network would be $\Delta L_{network} = \sum_{i=1}^{S} d = Sd$. The overall force that the network can exert would be:

$$F_{network} = \sum_{i=1}^{P} F = f(P, I_{ESA}, d, a, n, m, B_r, I, S) \quad (16)$$

which is a function of the following design parameters: number of fibers in parallel, length of a single ESA, distance between two coils of a single ESA, apothem of the polygonal profile of ESA, number of polygon’s side, coil density, residual flex density of the permanent magnet, the electric current and number of series ESAs within a fiber. All these parameters are real continuous variables except for number of fibers in parallel, number of polygon’s side, number of ESAs in a fiber and coil density that are discrete as well as residual flex density of the permanent magnet which is categorical variables. The goal is to maximize this nonlinear mixed-integer function as an objective function subjects to functional constraints regarding the overall size (length, width and thickness) that the network can occupy, the overall deflection of the network as well as its power consumption. These functional constraints depend on particular application of the network. For this work, the functionality of the networked ESAs will be examine as actuators for an active elbow brace.

III. OPTIMAL DESIGN BASED ON BRANCH AND BOUND ALGORITHM

In the previous section, the output force of a single ESA as well as that of a network was formulated as a function of design parameters. Obviously, the force function is nonlinear and as the design parameters are both real and categorical, the optimization process should be able to deal with nonlinear mixed-integer functions. Here we propose Branch and Bound algorithm to find the optimal solution to this problem.

Generally, the problem of designing machines (e.g. actuators) is understood and formulated as an inverse problem. The direct problem of the design can be defined as follows: From an actuator where the structure, the dimensions, and the materials are known, compute the performances; for example, the velocity, the torque, etc.

In the past years, the corresponding direct problem has first been studied by developing the well-known methods such as finite-element. Thus, with these numerical tools, the designer can iteratively generate and study some possible solutions of the considered inverse problem. Nevertheless, such a way cannot be entirely satisfactory because it does not solve the inverse problem rationally.
The corresponding inverse problem of design is: *From the values given by the desired performance (for example the torque), get the structure, the dimensions, and the material of the required actuator.* Such a problem is ill-posed in the Hadamard sense [27], because: 1) the existence and the uniqueness of the solution cannot be guaranteed and 2) this problem may generate an infinite number of solutions. Many counter-examples can be found, like the actuator considered in [28]. The inverse problem of design of electrical machines is explained in [29].

More recently, optimal dimensioning problems (that only deal with continuous variables) have been considered and solved by using classical local search algorithms, such as augmented Lagrangian method or Sequential Quadratic Programming (SQP) method, associated with analytical models [30], [31]. Many analytical models for different actuators have then been introduced depending on different assumptions [32]–[35]. Generally, in all these cases, only classical local search methods are considered and, therefore, only optimal dimensioning problems have been considered; the hypotheses used in this kind of algorithms can only deal with continuous optimization problems. Unfortunately, even for this kind of problem, it has been proven in [29] that such an inverse problem must be solved by using an exact global optimization algorithm and then it must be treated as a constrained-continuous global optimization problem.

In the classical literature about modeling and optimal design of actuators, some particular inverse problems, named predimensioning problems of design, were considered [36]. These works focus on the association of analytical models with local standard optimization algorithms [37]. Thus, first optimal solutions are found for some actuators. In this case, the solutions obtained by local optimization methods depend on the starting point introduced by the designer. Stochastic global optimization algorithms are not well adapted to solve this kind of problem because there are too many hard constraints. In [30], it is proven that even if a simple actuator is considered, local search solutions failed to find the optimal solution; see [36] to compare results with the exact ones published in [30].

In [32], [38], a rational way associating combinatorial analytical models of actuators and an exact global optimization algorithm, was proposed and studied. The algorithm was an exact global optimization method; see [30], [33], [35] for details on such algorithms and on their rigorous convergence to the global optimum. The obtained solutions satisfied the imposed set of conditions via combinatorial analytical models. When these optimal solutions were validated by the means of numerical tools (such as finite-element methods), some differences about the values of the electromagnetic torque were noted. Thus, the optimal results are not robust and this involves some adjustments of the parameters of the obtained solution. These adjustments can be done by iteratively solving the direct problem of design until the desired performance will be satisfied using a numerical model. In some numerical tools, such as ANSYS, stochastic and local optimization algorithms can be used to solve this problem of optimal adjustments.

The principle of this algorithms is based on subdivisions of the considered initial domain into smaller and smaller parts such that one can determine, using interval analysis [39], which box can be discarded. This is because interval computations will produce the proof that a box cannot contain the global solution or prove the fact that at least one constraint could not be satisfied in a box. Therefore, at the end of the algorithm, the global optimum will be enclosed with a given accuracy, however the algorithms must be capable to deal with both real continuous and categorical variables.

The main idea is to subdivide the initial domain space \( X \subseteq \mathbb{R}^n \times \prod_{k=1}^{n} K_k \) into smaller sub-boxes \( Z \subseteq X \) and to delete the considered box \( Z \), if and only if it can be proven that \( Z \) cannot contain the global optimum.

In order to subdivide the domain into smaller boxes, we need to consider the space in which the boxes are located in. In this problem, this space is a combination of continuous as well as categorical variables.

Branch and Bound algorithms have been modified to deal with integer variables, such as in our previous work in finding an optimal solution to the facility layout problem [40]. The categorical variables cannot be directly considered in subdivision of a box and an expression of a function. Here, we propose a method to change categorical variables into integer ones. This is done through an univariate function \( a \) that assigns an integer to a categorical value. For example, if one material has flux density of 0.85 Tesla and stiffness of another material’s flux density is 0.9 Tesla, this univariate function assigns integer value \( 1 \) to the first stiffness value and integer value \( 2 \) to the second one: \( a(1) = 0.85 \) Tesla and \( a(2) = 0.9 \) Tesla. Therefore, while integer values 1 and 2 can be used for the purpose of subdividing the boxes, the values \( a(1) \) and \( a(2) \) will be used in calculating the objective function and constraints.

The classical principle of subdividing is to choose a coordinate direction parallel to which \( Z \) has an edge of maximum length. Then, \( Z \) is subdivided normal to this direction [39]. For continuous variables the subdivision of a box \( Z \) on its kth component into to boxes \( Z_{1k} \) and \( Z_{2k} \) will be applied at the midpoint (as defined in [39]) of the original box \( Z \) as follow: \( Z_{1k} = \left[ Z_{k}^L + (Z_{k}^U + Z_{k}^L)/2 \right] \) and \( Z_{2k} = \left[ (Z_{k}^U + Z_{k}^L)/2, Z_{k}^U \right] \), where \( Z_{k}^L \) and \( Z_{k}^U \) denote the lower and upper bounds of kth component of box \( Z \), respectively. For categorical variables, however, this subdivision rule has to be slightly modified as follow: \( Z_{1k} = \left[ Z_{k}^L, [(Z_{k}^U + Z_{k}^L)/2] \right] \) and \( Z_{2k} = \left[ [(Z_{k}^U + Z_{k}^L)/2] + 1, Z_{k}^U \right] \), where \( \lfloor x \rfloor \) is the integer part of \( x \).

At each step of the Branch and Bound algorithm, and as the branches are generated through subdividing the boxes, we will have an interval for each variable. The interval analysis proposed by [39] is a powerful tool to calculate upper and lower bounds of a function over a box. This is simply done by the concept of *natural extension* of a function.
The natural extension of an expression of \( f \) into interval consists by replacing each occurrence of a variable by its corresponding interval (which encloses it), and then by applying the rules of interval arithmetic as explained in [39]. For example, if \( f(x) = x^2 + x + 1 \) and \( x \in X = [1, 2] \), then \( F(X) = ([1, 2])^2 + [1, 2] + 1 \) (\( X \) and \( F(X) \) are both intervals) is the natural extension of \( f(x) \) over box \( X \). It has been proven by [37] that the natural extension is always an inclusion function, which means that it encloses the upper and lower bounds of \( f \) over the box (or even all boxes when dealing with multiple variables) \( X \). Therefore, \( f(X) = [\min_x f(x), \max_x f(x)] \subseteq F(X) \forall X \).

Interval arithmetic is only defined for continuous functions, and thus in our method, the inclusion functions must be extended to deal with discrete variables as well. As mentioned before, we will convert categorical variables into integer variables through univariate functions. These integer variables will then be further relaxed into continuous variables. For example, if an integer variable belongs to set \( \mathbb{Z}^k \), \( \mathbb{Z}^{k+1} \), \( \mathbb{Z}^{k+2} \), \ldots, \( \mathbb{Z}^l \), a continuous interval \([\mathbb{Z}^k, \mathbb{Z}^l]\) will be considered for this variable. Replacing an integer set by its corresponding relaxed continuous interval, an inclusion function can then be constructed. The proof is obvious because it comes from the fact that the relaxed compact interval sets enclose by definition the initial discrete sets.

In order to proceed with the elimination of the boxes that cannot contain the global optimum solution, just the computation of lower bound of a given function \( f \) over a box \( Z \), i.e. \( lb(f, Z) \), is needed. We will compute these bounds by using interval analysis as follows: Considering that \( \tilde{f} \) denotes the current solution (in fact, it is just the best evaluation of \( f \) at this stage of the algorithm such that all the constraints are satisfied), one obtains the following:

1) No global solution is in \( Z \), if \( lb(f, Z) > \tilde{f} \), a lower bound of \( f \) over \( Z \) is greater than a solution already found, then no point in \( Z \) can be a global minimum.

2) No feasible solution is in \( Z \), if it exists \( k \) such that \( lb(g_k, Z) > 0 \), or it exists \( k \) such that \( lb(h_k, Z) > 0 \) or \( ub(h_k, Z) > 0 \) (\( ub = \) upper bound). In this case, a very small positive real value can be considered instead of 0, in order to address numerical approximation difficulties due to the floating point operations [39].

This implies that at the end of the algorithm, accurate enclosures of the global minimum value and of all their corresponding solutions can be expected. Indeed, at each step of the algorithm, one has the following properties for all the remaining sub-boxes \( Z \subseteq X \): 1) \( lb(f, Z) \leq \tilde{f} \) and 2) all constraints are always satisfied.

### IV. CASE STUDY: AN ACTIVE ELBOW BRACE

As mentioned before, bio-inspired arrangement of ESAs can lead to an enhanced output force. If we can achieve the level of forces that are exerted physiologically at a human’s extremity joints, by on-board batteries, this actuation technology can greatly impact rehabilitation or force augmentation applications for mobility impaired patients, as it is light weight, has high bandwidth, and it is powerful and portable.

To determine whether or not networked ESAs has this capability, we selected a human elbow joint, as a critically important joint in performing our daily activities. The elbow joint functions as a functional link between the upper arm and the forearm. It positions the hand in space, allows the forearm to act as a lever in lifting and carrying, and provides precision in both open and closed kinetic chain work. Therefore, even a mild impairment in elbow can significantly reduce the ability of the hand to reach its objectives, with the problem being compounded by the fact that the other joints in the upper limb are unable to compensate for this loss.

#### A. CONSTRAINTS

In this section, we first determine the performance requirement of an active elbow brace that is supposed to be worn by a patient and match the performance of a healthy elbow joint [41], in terms of maximum torque and angular deflection, i.e. flexion-extension. Range of elbow’s flexion-extension in a healthy individual is between 0 to 104 degree, while the maximum torque, that the elbow can handle at 90 degree (flexed) is about 4-6 Nm [42]. The other performance parameters of the active elbow brace that are summarized in Table 1, have been calculated based on available passive elbow braces in the market.

We consider a popular elbow brace by Orthomen [43]. To make this passive brace active, two networked ESAs can be attached to the brace in an agonistic-antagonistic fashion, similar to the Biceps and Triceps skeletal muscles. Hence,

---

**Algorithm 1 Branch and Bound Algorithm**

1. set initial box \( X := \mathbb{R}^n \times \prod_{i=1}^n K_i \)
2. set current solution \( \tilde{f} = +\infty \)
3. set initial lower bound \( \ell = (+\infty, \mathbb{X}) \)
4. extract from \( \ell \) the lowest lower bound \( lb \)
5. divide the chosen box \( V \) by its midpoint \( m \) to get sub-boxes \( V_1 \) and \( V_2 \)
6. for \( j = 1, 2 \)
   a. calculate \( v_j = lb(f, V_j) \)
   b. calculate lower bounds of all constraints over \( V_j \) using interval analysis
   c. if \( \tilde{f} \geq v_j \) and all constraints are satisfied then:
   i. set \( \ell := (v_j, \mathbb{V}_j) \)
   ii. if all constraints are satisfied at \( m \)
   iii. if \( \tilde{f} \) has been changed from its previous value, then remove all \((z, \mathbb{Z})\) from \( \ell \) where \( z > \tilde{f} \)
   7. if \( \min_{z, \mathbb{Z}} z \leftarrow \tilde{f} \) then STOP
   else GoTo step 4
8. end if
9. end for
the thickness limit of the network. We limit the maximum number of sides of the cross-section polygon to 8, however to consider circular cross-section profile, we would also assign a large number (e.g. 100) as for the polygon’s number of sides. Coil loop density is limited to 5000 loop per meter, due to the smallest feasible diameter size of the flexible PDMS wire (i.e. 0.1mm). The categorial variable, i.e. flux charge density has to be experimentally quantified. As mentioned before, to prepare the flexible permanent magnet, a mixture of PDMS and magnetic particles were cured inside a 3D printed mold while the whole mixture was exposed to a strong external magnetic field. However, the ratio between magnetic particles and PDMS plays a very important interdisciplinary role; i.e. both in magnetic and mechanical domain. By increasing the ratio of magnetic particles, the flux charge density increases that would lead to enhancement of the magnetic field and the force generated by the coils. However, in mechanical domain, this increase in the magnetic particle ratio, affects the Young Modulus (i.e. elasticity) of the flexible magnet as well as its yield point, which in turn, affects the maximum passive elongation ratio of the permanent magnet. We prepared 5 flexible permanent magnet samples each with different mixing ratios (between 8% to 28% with incremental mixing percentage of 4% of magnetic particles). For each sample, the flux charge density was measured by Magnetic Field Instrument (MFI), a device used to measure the magnetic field or flux around permanent magnets, coils, and electrical devices. The flux density for each mixing ratio was presented in an interval domain as the mixing ratio could not be precisely adjusted.

The result is shown in Table 2. The optimization problem then can be formulated as:

\[
\begin{align*}
\text{Max} & \quad F_{\text{network}} = \sum_{i=1}^{P} F = f(P_c, P_r) \\
\text{subject to:} & \\
L_{\text{network}}(P_c, P_r) & \leq 0.24m, \\
\Delta L_{\text{network}}(P_c, P_r) & \geq 0.09m, \\
T_{\text{network}}(P_c, P_r) & \leq 0.02m, \\
W_{\text{network}}(P_c, P_r) & \leq 0.05m, \\
\vec{P}_r & = [I_{\text{ESA}}, a, d, I]^T \\
\vec{P}_c & = [m, n, B_r, P, S]^T \\
I_{\text{ESA}} & \leq 0.24m \\
a & \leq 0.02m \\
d & \leq 0.09 \\
I & \leq 10\text{Amp} \\
n & \in [1, ..., 5000] \\
P & \in [1, ..., 100] \\
S & \in [1, ..., 100] \\
B_r & \in [1, ..., 5] : \text{Table 2} \\
m & \in [1, ..., 8] \text{ and } 100
\end{align*}
\]

where \( P_c \in IR^8 \) and \( P_r \in \prod_{i=1}^{5} \) are real and discrete variable vectors, respectively. With these settings, we performed the proposed Branch and Bound optimization process to find

| Parameter | Value | Unit |
|-----------|-------|------|
| Range of Flexion-Extension | 0-104 | Degree |
| Weight | 150-300 | g |
| Volume | 200 | cm³ |

**TABLE 1. Performance of active elbow brace.**
what design parameters would lead to maximum output force and determined if that optimal design parameters can satisfy the aforementioned performance and input parameters.

C. OPTIMIZATION RESULTS

We conducted mechanical tensile tests on several mixtures of PDMS and magnetic particles with different mixing ratios. The results suggested that mixing ratio of magnetic up to 28% (with respect to the weight compared to that of PDMS) would lead to safe (i.e. passive elongation ratio > 35%).

Result of the optimization problem is presented in Table 3.

| Parameter                          | Value | Unit |
|------------------------------------|-------|------|
| coil density                       | 5000  | loop/meter |
| apothem                            | 2.5   | mm    |
| length of ESA                      | 26    | mm    |
| distance between coils             | 5     | mm    |
| number of polygon sides            | 6     |      |
| electric current                   | 10    | Amp   |
| number of series ESAs within a fiber | 14   |      |
| number of parallel fibers within a network | 20-24 | %      |

The Results on the optimal cross-section profile showed that although for a stacked network of ESAs, the triangular cross-section profile makes the largest magnetic field and force at the center of coil compared to the other profiles, it also consumes the most electrical power such that after normalizing and obtaining the ratio, the triangular profile is the least efficient section profile while the circular one is the most.

We also considered the effect of packing density of the profiles since all of the coil’s section profiles do not cover the available space entirely and there are some free spaces among the profiles for circular, octagonal and pentagonal arrangements. For this purpose the packing density coefficient ($\lambda = \frac{P_{A_{coil-section}}}{A_{network}}$) was defined and calculated for each section profiles. After applying the packing density coefficient to the previous findings we found that the hexagonal arrangement for the coil’s section profile ($m = 6$) is the most efficient one due to have the largest amount ($\lambda = 1$) among all section profiles.

Furthermore, the optimal agnostic network on each side (left and right) should consist of 40 fibers, attached in parallel, each has of 9 ESAs in series for the agnostic networks. With these networks the total weight of the active brace would be 200 g, while the volume stays below 200 $cm^3$ and the flexion-extension range achieves 93 degrees rotation around the elbow joint (11 degrees less that what is actually required). The force of the agnostic network as a result of optimization is 71 N which means that the output torque of the active brace would be 6 Nm that matches the required output torque at the elbow according to [42].

In order to evaluate the result of the optimization process, we have built a small network based on the optimal design parameters with 3 parallel fibers where each fiber consists of 4 series ESAs as it is shown in 9. The small network was able to lift a 2.5N weight, with its 3 fibers. When scaled up to the number of fibers of agnostic network suggested by the optimization process, the total output force would be 66.6N, which is slightly less than what was predicted by the optimization process (i.e. 71N). A reason could be due to actual coil density. As it is clear from Fig. 9, there are gaps between the micro-channels inside an ESA, as the flexible wire was manually wrapped around the coil. This would lead to less number of the loops per meter and eventually less output force. An automated wrapping technique would solve this issue.

V. CONCLUSION AND FUTURE WORKS

In this paper, optimal design for a network of novel ESA based on branch and bound algorithm was presented. The novel electromagnetic soft actuator operates based on the working principle of Solenoids, that consists of two antagonistically located coils made of flexible wires and share a flexible permanent magnetic core.

It was shown that by reducing the size of these ESAs, the force to volume size ratio increases, which suggest a network of miniaturized ESAs would achieve higher amount of output force compare to a single ESA with the same size of the whole network. In this work this was numerically tested for a case study of an active elbow brace.
Branch and Bound optimization algorithm was employed to achieve optima design of a single ESA and consequently optimal design of a networked ESAs to achieve the maximum output torque at the elbow joint, while the performance parameters are being satisfied.

The result showed that having a network of ESAs as drive train for an active brace, we can satisfy the performance parameters, for supporting the elbow joint of a patient with decreased muscle performance and mobility.

This suggests that with the available manufacturing process discussed in this paper, the actuation technology based on electromagnetic soft actuators can to be used as drive trains in robotic prosthesis and robotic exoskeletons, to support patients with decreased muscle function at their affected joints. Our future endeavors are focused in enhancing further the produced torque by the ESAs, in order to be utilized for robotic prosthetics and exoskeletons in patients with complete loss of muscle function.

This actuation technology is uniquely suitable in rehabilitation and/or force augmentation applications for those mobility impaired patients that have not completely lost the ability to move their affected joints and would need some extra help to recover or be able to perform their daily tasks. Considering the huge population of these types of mobility impaired patients (e.g stroke patients, peripheral arterial disease, traumatic injuries, neuropathies, senescence and frailty) electromagnetic soft actuators provide novel potential solutions, for wearable and next-to-skin type of assistive technologies, at low production cost, safe, portable and yet sufficiently powerful with low power requirement and high bandwidth.

Planning for the future, we will manufacture the active brace powered with the network of ESAs as suggested by the Branch and Bound algorithm and test its potentials in rehabilitation trials, such as those for elbow stiffness or iso-kinetic motion for elbow spasticity and other motor dysfunctions from various medical pathologies.

REFERENCES

[1] S. Haddadin, A. Albu-Schäffer, and G. Hirzinger, “Safety evaluation of physical human-robot interaction via crash-testing,” in Robotics: Science and Systems, vol. 3. Cambridge, MA, USA: MIT Press, Jun. 2007, pp. 217–224.

[2] A. De Santis, B. Siciliano, A. De Luca, and A. Bicchi, “An atlas of physical human–robot interaction,” Mech. Mach. Theory, vol. 43, no. 3, pp. 253–270, Mar. 2008.

[3] S. Haddadin, A. Albu-Schäffer, and G. Hirzinger, “Requirements for safe robots: Measurements, analysis and new insights,” Int. J. Robot. Res., vol. 28, nos. 11–12, pp. 1507–1527, Nov. 2009.

[4] R. Alami, A. Albu-Schaeffer, A. Bicchi, R. Bischoff, R. Chatila, A. De Luca, A. De Santis, G. Giralt, J. Guiochet, G. Hirzinger, F. Ingrand, V. Lippiello, R. Mattone, D. Powell, S. Sen, B. Siciliano, G. Tonietti, and L. Villani, “Safe and dependable physical human-robot interaction in anthropic domains: State of the art and challenges,” in Proc. IEEE/RSJ Int. Conf. Intell. Robots Syst., Oct. 2006, pp. 1–16.

[5] C. Lee, M. Kim, Y. J. Kim, N. Hong, S. Ryu, H. J. Kim, and S. Kim, “Soft robot review,” Int. J. Control, Autom. Syst., vol. 15, no. 1, pp. 3–15, Feb. 2017.

[6] K. Suzumori, S. Endo, T. Kanda, N. Kato, and H. Suzuki, “A bending pneumatic rubber actuator realizing soft-bodied manta swimming robot,” in Proc. IEEE Int. Conf. Robot. Autom., Apr. 2007, pp. 4975–4980.

[7] P. Polygerinos, S. Lyne, Z. Wang, L. F. Nicolini, B. Mosadeq, G. M. Whitesides, and C. J. Walsh, “Towards a soft pneumatic glove for hand rehabilitation,” in Proc. IEEE/RSJ Int. Conf. Intell. Robots Syst., Nov. 2013, pp. 1512–1517.

[8] H. K. Yap, J. Hoon Lim, F. Nasrallah, J. C. H. Goh, and R. C. H. Yeow, “A soft exoskeleton for hand-assistive and rehabilitation application using pneumatics actuators with variable stiffness,” in Proc. IEEE Int. Conf. Robot. Autom. (ICRA), May 2015, pp. 4967–4972.

[9] A. Mata, A. J. Fleischman, and S. Roy, “Characterization of polydimethylsiloxane (PDMS) properties for biomedical Micro/Nanosystems,” Biomed. Microdevices, vol. 7, no. 4, pp. 281–293, Dec. 2005.

[10] H. R. Choi, S. M. Ryu, K. M. Jung, H. M. Kim, J. W. Jeon, J. D. Nam, R. Maeda, and K. Tanie, “Microrobot actuated by soft actuators based on dielectric elastomer,” in Proc. IEEE/RSJ Int. Conf. Intell. Robots Syst., vol. 2, 2002, pp. 1730–1735.

[11] G.-Y. Gu, U. Gupta, J. Zhu, L.-M. Zhu, and X. Zhu, “Modeling of viscoelastic electromechanical behavior in a soft dielectric actuator motor,” IEEE Trans. Robot., vol. 33, no. 5, pp. 1263–1271, Oct. 2017.

[12] C.-T. Nguyen, H. Phung, T. D. Nguyen, H. Jung, and H. R. Choi, “Multiple-degrees-of-freedom dielectric elastomer actuators for soft-printable hexapod robot,” Sens. Actuators A, Phys., vol. 267, pp. 505–516, Nov. 2017.

[13] M. Duduta, E. Hajesmaili, H. Zhao, R. J. Wood, and D. R. Clarke, “Realizing the potential of dielectric elastomer artificial muscles,” Proc. Nat. Acad. Sci. USA, vol. 116, no. 7, pp. 2476–2481, Feb. 2019.

[14] H. Rodrigue, W. Wang, D.-R. Kim, and S.-H. Ahn, “Curved shape memory alloy-based soft actuators and application to soft gripper,” Composite Struct., vol. 176, pp. 398–406, Sep. 2017.

[15] L. Fang, T. Fang, X. Liu, S. Chen, C. Lu, and Z. Xu, “Near-infrared light triggered soft actuators in aqueous media prepared from shape-memory polymer composites,” Macromol. Mater. Eng., vol. 301, no. 9, pp. 1111–1120, Sep. 2016.

[16] A. Miriyev, K. Stack, and H. Lipson, “Soft material for soft actuators,” Nature Commun., vol. 8, no. 1, pp. 1–8, Dec. 2017.

[17] H. Lipson, “Challenges and opportunities for design, simulation, and fabrication of soft robots,” Soft Robot., vol. 1, no. 1, pp. 21–27, Mar. 2014.

[18] T. N. Do, H. Phan, T.-Q. Nguyen, and Y. Visell, “Miniature soft electromagnetic actuators for robotic applications,” Adv. Funct. Mater., vol. 28, no. 18, May 2018, Art. no. 180024.

[19] Y.-P. Yang, J.-J. Liu, D.-H. Ye, Y.-R. Chen, and P.-H. Lu, “Multiobjective optimal design and soft landing control of an electromagnetic valve actuator for a camless engine,” IEEE/ASME Trans. Mechatronics, vol. 18, no. 3, pp. 963–972, Jun. 2013.

[20] T. Hellebrekers, O. Kromer, and C. Majidi, “Soft magnetic skin for continuous deformation sensing,” Adv. Intell. Syst., vol. 1, no. 4, Aug. 2019, Art. no. 1900025.

[21] W. Bartlett and R. J. Wood, “Comparative analysis of fabrication methods for achieving rounded microchannels in PDMS,” J. Micromech. Microeng., vol. 26, no. 11, Nov. 2016, Art. no. 115013.

[22] N. Ebrahimi, S. Nugroho, A. F. Taha, N. Gatisis, W. Gao, and A. Jafari, “Dynamic actuator selection and robust state-feedback control of networked soft actuators,” in Proc. IEEE Int. Conf. Robot. Autom. (ICRA), May 2018, pp. 2857–2864.

[23] A. Jafari and N. Ebrahimi, “Electromagnetic soft actuators,” U.S. Patent App. 16/457,452, Jan. 2, 2020.

[24] N. Ebrahimi, P. Schimpf, and A. Jafari, “Design optimization of a solenoid-based electromagnetic soft actuator with permanent magnet core,” Sens. Actuators A, Phys., vol. 284, pp. 276–285, Dec. 2018.

[25] B. A. Kemp, T. M. Grzegorczyk, and J. A. Kong, “Lorentz force on dielectric and magnetic particles,” J. Electromagn. Waves Appl., vol. 28, no. 8, pp. 827–839, Jun. 2014.

[26] P. T. Pappas, “The original amperes force and biot-savart and lorentz forces,” Il Nuovo Cimento B, vol. 76, no. 2, pp. 189–197, Aug. 1983.

[27] M. M. Lavrentiev, V. G. Romanov, and S. P. Shishatski, Ill-Posed Problems of Mathematical Physics and Analysis, vol. 64. Providence, RI, USA: American Mathematical Society, 1986.

[28] E. Fitan, F. Mesine, and B. Nogarede, “The electromagnetic actuator design problem: A general and rational approach,” IEEE Trans. Magn., vol. 40, no. 3, pp. 1579–1590, May 2004.

[29] F. Mesine, B. Nogarede, and J.-L. Lagouanelle, “Optimal design of electromechanical actuators: A new method based on global optimization,” IEEE Trans. Magn., vol. 34, no. 1, pp. 299–308, 1st Quart., 1998.

[30] A. D. Kone, B. Nogarede, and M. Lajoie Mazenc, “Le dimensionnement des actionneurs électriques: Un problème de programmation non linéaire,” J. Phys. I France, vol. 3, no. 2, pp. 285–301, Feb. 1993.
[31] F. Wurtz, J. Bigeon, and C. Poirson, “A methodology and a tool for the computer aided design with constraints of electrical devices,” IEEE Trans. Magn., vol. 32, no. 3, pp. 1429–1432, May 1996.

[32] E. Fitan, F. Messine, and B. Nogarede, “A general analytical model of electrical permanent magnet machine dedicated to optimal design,” COMPET. Int. J. Comput. Math. Electr. Electron. Eng., vol. 22, no. 4, pp. 1037–1050, Dec. 2003.

[33] S. Huang, M. Aydin, and T. Lipo, “A direct approach to electrical machine performance evaluation: Torque density assessment and sizing optimization,” in Proc. Int. Conf. Electr. Mach., 2002, Art. no. 235.

[34] T. Miller, Brushless Permanent-magnet and Reluctance Motor Drives, Oxford, U.K.: Oxford Science Publications, 1989.

[35] G. R. Slemon and X. Liu, “Optimal solution for the two-dimensional facility layout problem using a branch-and-bound algorithm,” Comput. Ind. Eng., vol. 55, no. 3, pp. 606–619, Oct. 2008.

[36] C. D. Bryce and A. D. Armstrong, “Anatomy and biomechanics of the elbow,” Orthopedic Clinics North Amer., vol. 39, no. 2, pp. 141–154, 2008.

[37] R. Hosoda and G. Venture, “Human elbow joint torque estimation during dynamic movements with moment arm compensation method,” IFAC Proc. Volumes, vol. 47, no. 3, pp. 12305–12310, 2014.

[38] Orthomen. Post Op Elbow Brace. Accessed: 2007. [Online]. Available: https://www.orthomen.com/

[39] NAFISEH EBRAHIMI received the bachelor’s degree in mechanical engineering from the Isfahan University of Technology (IUT), Iran, in 2007, and the master’s degree in mechanical engineering from the Iran University of Science and Technology (IUST), Iran, in 2010. She is currently pursuing the Ph.D. degree with the Advanced Robotic Manipulators (ARM) Laboratory, Department of Mechanical Engineering, College of Engineering, The University of Texas at San Antonio (UTSA). Her research interests include soft robotics, rehabilitation robotics, mechatronics, and biomechanics.

[40] TEJA GUIDA received the B.Tech. degree from the Indian Institute of Technology, Bombay, India, in 2004, and the Ph.D. degree in biomechanical engineering from The University of Texas at San Antonio / University of Texas Health Science Center at San Antonio, in 2008. From 2008 to 2014, he was a Postdoctoral Researcher with the Wake Forest Institute for Regenerative Medicine. His interests are focused on developing regenerative strategies for bone and skeletal muscle tissue engineering. In specific, on-going projects focus on the employment of mechanical stimulation regimes in bioreactors to improve engraftment of bone and muscle substrates. The compliance of substrates has a profound influence on the stem cell response and this effect is enhanced by applied mechanical loads. Studying the relative impact of these cues on directed cellular differentiation is important to develop appropriate materials for musculoskeletal interfaces, where there is a steep gradient in mechanical properties. This work is being carried out in close collaboration with Dr. Joseph Wenke’s group at the US Army Institute of Surgical Research. Dr. Guda was a recipient of the National Science Foundation NSF-CAREER Award, in 2019.

[41] MILTIADES ALAMANIOTIS (Member, IEEE) received the Diploma degree in electrical and computer engineering from the University of Thessaly, and the Ph.D. degree in nuclear engineering emphasizing on applied artificial intelligence from Purdue University, in 2012. He has been an Assistant Professor with the Department of Electrical and Computer Engineering, The University of Texas at San Antonio (UTSA), since August 2018. Prior to joining UTSA, he worked as a Research Assistant Professor at Purdue University. His interdisciplinary research focuses on the development of artificial intelligence approaches applied to energy and electrical power systems, smart cities, nuclear energy systems, and nuclear security. His main research interests focus on applied artificial intelligence in various engineering applications, such as smart grids, smart cities, nuclear power, and nuclear security.

[42] AMIR JAFARI (Member, IEEE) received the Ph.D. degree in robotics from the Italian Institute of Technology (IIT), Genoa, Italy, in 2011. From 2011 to 2012, he was a Postdoctoral Researcher with IIT. From 2012 to 2013, he was a Postdoctoral Researcher with ETH in Zurich, Zürich, Switzerland. From 2014 to 2016, he was with the Agency of Science, Technology and Research, Singapore, before he joined The University of Texas at San Antonio (UTSA), San Antonio, TX, USA, in 2016. He is currently an Assistant Professor with the Department of Mechanical Engineering, UTSA. His research areas include physical human–robot interaction, compliant/variable compliance mechanisms, and rehabilitation robotics.

[43] DIMITRIOS MISERLIS (Member, IEEE) received the M.D. degree from the University of Crete, Heraklion, Greece, in 2008. From 2009 to 2018, he was with the University of Nebraska Medical Center, Omaha, NE, USA, as a Postdoctoral Research Fellow in vascular surgery biomedical research, as well as a General Surgery Residency training and a Vascular Surgery clinical Fellowship. He is currently a Vascular Surgeon with the University of Texas Health at San Antonio (UTHSA) and an Assistant Professor with the Biomedical Engineering Joint Graduate Program, The University of Texas at San Antonio (UTSA). His research areas include vascular surgery biomedical technologies, imaging, flexible robotics, and artificial intelligence integration.