Guided-Wave Up-Conversion Luminescence in Er\(^{3+}/\)Yb\(^{3+}\) Co-doped Phosphate Glass Waveguide Produced by Direct Femtosecond Laser Writing

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We report on guided-wave up-conversion luminescence in femtosecond laser writing cladding waveguides in Er\(^{3+}/\)Yb\(^{3+}\) co-doped phosphate glass. The waveguides were fabricated with 30-\(\mu\)m- and 100-\(\mu\)m-diameter of the guiding cores. The guiding properties of waveguides have been investigated at 633 nm by end face coupling of free space light and physical contact of fibers. The experimental and calculated results of propagating modal profiles and losses have proved favorable performances suitable to Gaussian mode field and multi-mode applications. Under the optical pump laser at 980 nm, the guided-wave up-conversion luminescence at visible light range has been realized through the waveguides.

Keywords: optical waveguides, rare-earth-ion doped glass materials, up-conversion, femtosecond laser micromachining, integrated optical device

INTRODUCTION

Rare-earth-doped glasses are widely applied in optical communications, optical sensing, imaging, etc., benefiting for their outstanding material properties including high ions solubility, prominent nonlinear effects, and adjustable refractive indices [1, 2]. Many kinds of glasses, such as silicate glass, chalcogenide glass, phosphate glass, tellurite glass, and fluoride glass were successfully used to fabricate active integrated devices and lasers, by doping rare-earth ions [3–5]. In these glasses, the silicate as well as phosphate glasses are most commonly used for preparing fibers operating around 1 \(\mu\)m. The phosphate glass can realize a high concentration up to above 20 wt%, contributing to extremely enhance the pump absorption efficiency and thus shorten the fiber length. Thus far, Nd\(^{3+}\) ion doped, Er\(^{3+}\) ion glass and Yb\(^{3+}\) ion doped phosphate glass have been developed in fields of laser fusion, laser weapon, laser ranging, optical communication waveguide amplifier and ultrashort pulse laser [6, 7]. Er\(^{3+}\) ion has been regarded as the most efficient candidate to generate up-conversion luminescence due to its appropriate electronic structures [8]. The optical materials co-doped with erbium (Er\(^{3+}\)) and ytterbium ions (Yb\(^{3+}\)) exist energy transfer between Er\(^{3+}\) and Yb\(^{3+}\), which provides a straightforward way to combine the Yb\(^{3+}\) ion’s simple energy level scheme with the Er\(^{3+}\) ion’s electronic states, for low-cost pumping at 980 nm [9]. Due to a few different mechanisms of Er\(^{3+}\) ions including exited-state absorption and the photon avalanche effect, these materials can be excited the up-conversion luminescence that the wavelength of emission light is shorter than the excitation. Currently, there are great interests in luminous materials for efficient up-conversion from near
infrared to visible emission, which is applied in extensively fields of super-resolution imaging, medical diagnoses, sensors [10–12]. Phosphate glasses, mainly composed of P2O5 instead of SiO2 in silica glasses, can be beneficial over silica glasses with low transform temperature, making the materials to be suffered from heat induced end-face damage at a relative high output power, commonly employed for optical integrations [13, 14].

Waveguides, as a fundamental component in optical integrated circuits, have been achieved for a number of photonic applications in rare-earth-ion doped materials. Benefitting from compact geometry of waveguides and combination of versatile bulk features, miniaturized devices on waveguide platforms can be fabricated with high density, low optical thresholds, on-chip sizes, etc. Optical waveguide is able to localize the propagation of light field within the space at micron scale, thus increase the light density considerably with respect to that in the bulk material. Optical waveguides can be fabricated in many ways, such as ion implantation proton exchange, thermal ion diffusion, thin film deposition, femtosecond laser processing, etc., [15–19]. As a new technology of fabricating waveguide, femtosecond laser can be focused in the material to change the refractive index distribution of the material through a few nonlinear effects to form a waveguide structure [20–23]. Femtosecond laser is often used to prepare optical waveguides in rare-earth-doped materials. In recent years, several phenomenon and applications have been discovered by using femtosecond laser for optical waveguide manufacturing in rare-earth-doped glass materials, including planar and channel waveguides [24–26]. In 2016, D.G. Lancaster et al. reported an Er3+ Yb3+ co-doped waveguide laser fabricated by fs laser writing [27]. According to refractive index changes of writing track in all optical materials, laser-writing configuration can be categorized into positive change and negative change. Different from positive changes (so-called Type I with directly refractive index increasing, Type II with laser induced stress change), negative changes are often used to make depressed cladding waveguides (Type III). This type of waveguide can be arbitrarily controlled its cross-section size in the range of ten microns to hundreds of microns [28–30]. In 2018, S.N.C. Santos et al demonstrated the fabrication of 3D Type I waveguides in barium borate glass co-doping with Dy3+ and Eu3+ ions for the visible band applications [31]. In 2021, M. Khalid et al showed the femtosecond laser writer Type II waveguides in Yb doped glass for 1 μm laser applications [32]. Compared with Type I and II, depressed cladding waveguides commonly are considered to possess better performance in fiber-to-waveguide coupling with adjustable diameter of the end-face cross-section [15, 28–30].

In this work, we have achieved for depressed cladding waveguides with two different cross-section sizes by femtosecond laser writing in Er3+ and Yb3+ co-doped phosphate glass. The up-conversion luminescence of glass waveguides under visible light was studied under NIR optical pumping. Through the end coupling system, we study the waveguide characteristics, transmission mode and transmission loss of the waveguide at 633 nm. Under pump laser at 980 nm, we successfully realized the guided-wave up-conversion luminescence at visible light range through the waveguides and discussed the possible up-conversion mechanism based on the energy match situation dependency on pumping power and the Er3+ and Yb3+ ion energy level diagram.

EXPERIMENTS IN DETAILS

The phosphate glass for waveguide fabrication is a mature commercial glass made by Shanghai Institute of Optics and Fine Mechanics, Chinese Academy of Science, called WM13. The glass has compositions: P2O5-Na2O-BaO-Al2O3-Er2O3-Yb2O3, with 1.5wt.% Er2O3 and 4.5wt.% Yb2O3, which has been cut into dimensions of 10 (x) 10 (y) × 2 (z) mm3 and optically polished. The cladding waveguides with circular boundaries were produced by using a home-made laser fabricating facility. The femtosecond laser system (amplified Ti: Sapphire, Spitfire, Spectra Physics, United States), with 120 fs duration, 800 nm central wavelength, and 1 kHz repetition rate, was used to generate linearly polarized laser to fabricate waveguides. The laser irradiation was controlled with a mechanical shutter. To control the shape of the beam and the energy on sample, a neutral density filter was placed before the focusing objective lens. The ×40 focusing microscope objective (Numerical Aperture N.A. = 0.40) was used to focus the laser beam at the depth of ~150 μm. And the phosphate glass sample was inserted on the XYZ translation stage with a constant moving velocity of 0.5 mm/s (the schematic plot shown in Figure 1A). The procedure was repeated at different depths and positions of the sample following the desired circular geometry. As result, two tubular structures containing a chain of parallel tracks were inscribed in the phosphate crystal by the fs-laser. The two depressed cladding waveguides were produced to be 30- and 100-μm-diameter with a pulse energy of ~0.42 μJ, shown in Figures 1B, C. The careful choice of the irradiation parameters (scanning velocity and pulse energy) avoided the formation of cracks inside the crystal and allowed the generation of cladding waveguides.

The guidance measurement was performed through an end-face coupling platform. The components in the platform

![Figure 1](https://example.com/figure1.png)
are adjusted when testing waveguide performance and guided-wave luminescence. For guidance performance, a linearly polarized He-Ne laser at wavelength of 633 nm was utilized as the light source, as shown in Figure 2. Placing a half-wave plate behind the laser converted the linearly polarized direction. The incident linearly polarized light was into the waveguides via a $\times 20$ microscope objective lens. Another microscope objective lens was used to out-couple the light, and finally a CCD camera as a detector collected and imaged the light (a photodetector altered to measure the output power). Through the end coupling system, we study the propagation modes and propagation losses of the waveguide at 633 nm. The propagation losses of the photonic structures were determined by directly measuring the light powers (by power meters) from the input and output end-faces. The propagation loss can be estimated with the following equation [33].

$$P_{out} = P_{in}(1 - R)^2 e^{-\alpha L T}$$

where $R$ is the Fresnel reflection coefficient. $L$ stands for the length of the waveguide. $T$ is related to the mismatch between the pump beam mode and waveguide mode. The value of mismatch is calculated by a numerical calculating software to be $\sim 19$ and $23\%$ for 30 nm- and 100 nm-diameter waveguides, respectively (more details discussed in next section).

For guided-wave luminescence, a filter was inset the platform behind collected objective for removing optical pumping signals. A spectrometer (Omni-λ750, Zolix Instruments Co. Ltd.) was used as the detector in the platform for measuring the spectra of the up-conversion at room temperature.

**RESULTS AND DISCUSSION**

Figures 3A, B shows the optical microscope images of the light entering the bulk and the linearly polarized waveguide of phosphate glass, respectively. The up-conversion luminescence light fast lost in the bulk interior. Therefore, a bright and short beam is observed at the interface of the bulk and fiber, which excited by 980 nm light pumping. On the contrary, the beam can be well propagated along the optical waveguide, and the light intensities are higher in the waveguide cores than in the bulk (the intensities can be detected in the waveguides, but not in the bulk). It means that the up-conversion luminescence light can be effectively limited in the waveguide written by the femtosecond laser.

Figure 4 depict the measured spatial modal profiles at the output of waveguide as the linearly polarized light at 633 nm is launched into the waveguides of 30- and 100-μm-diameter. An additional linearly polarizer is inserted into the end-face coupling system (before the detector, not shown in Figure 2) to inspect the polarization features of the output beams. It is worthy to point out that 30- and 100-μm-diameter waveguides could support arbitrary polarizations, and the mode distributions do not exhibit significant polarization dependence as a result of the symmetric morphology of the cladding structures. The TE (TM) modal profiles are demonstrated in Figures 4A,C (Figure 4B,D) with the diameters of 30 and 100 μm, respectively. Propagating beams supports near Gaussian mode and multi-mode excitation in 30- and 100-μm-diameter waveguides. And the polarizations of the propagating beams are well maintained in the both core-size waveguides.
By assumption of a step-like refractive index profile and measuring the numerical aperture of the waveguides, we can estimate the maximum refractive index change of the waveguide core using the formula [34]:

$$\Delta n = \frac{\sin^2 \theta_m}{2n}$$

where $\theta_m$ is the maximum incident angular deflection at which no transmitted power change occurs. All the measurements were carried out at room temperature. $n = 1.60$ is the refractive index of the bulk at 633 nm wavelength. By inserting a prism in the end-face coupling system, the incident angle ($\theta$) can be changed. The maximum incident angular ($\theta_m$) can be confirmed to be $\sim 3.5'$ with a difference of $0.3'$ between both polarizations. In this work, we calculated maximum refractive index change ($\Delta n$) $0.3 \times 10^{-3}$. The refractive index at the end-face of 30-μm-diameter waveguide has been reconstructed (as shown in Figure 5A). The modal profiles are calculated by the FD-BPM algorithm (Rsoft® Beam PROP), which are shown in Figures 5B,C. The mismatch of the launched Gaussian beam ($E_g$) with modal profiles ($E_p$) of waveguides can be calculated by an overlap integral between both profiles. The mismatch coefficient $L$ can be expressed as [35]:

$$T = \frac{1}{2} \left( \frac{E_p}{E_g} \right) \int S E_g \cdot E_p dS$$

The calculated data are in good agreements with the experimental ones. It could be clearly seen that the two waveguides fabricated with the same parameters except size also exhibit the positively related refractive index alternations.

Based on the end-face coupling system at 633 nm, we measured the light power of the input and output end-faces. The propagation losses were calculated to be $1.8 \pm 0.3$ dB/cm and $2.3 \pm 0.3$ dB/cm for 30- and 100-μm-diameter waveguide. It indicates that, the larger the diameter of the waveguide brings the bigger the propagation loss due to some effects introduced by multi-mode, for instance mode-coupling. Consequently, the propagation losses can be reduced by optimizing the size of the optical waveguide (e.g., reducing the diameter of the waveguide) of the femtosecond laser writing [36, 37].

Up-conversion luminescence is well-known process, which converts near-infrared radiation (NIR) into visible light. An anti-Stokes process can convert two or more low energy photons into a higher energy photon. Here, a 980-nm continuous-wave laser was used to excite up-conversion luminescence through an end-face coupling system of optical fiber-waveguide physical contact end-surfaces. As shown in Figure 6A, emission spectra were collected by a micro-objective output from the waveguides and detected by a spectrometer in the 400–700 nm range with a resolution of 0.1 nm. The spectra of bulk and waveguide have two peak positions at 650 and 660 nm, and the peak positions are almost the same. From 400 to 560 nm, there is almost no difference between the emission spectra of bulk and waveguide. From 560 nm, the emission spectrum of the waveguide began to increase significantly. The emission spectra are higher in the waveguides than in the bulk. For the up-conversion emissions, the intracavity light intensities from the waveguides are increased by 38.3% in comparison to the bulk. There is little difference between waveguides with diameters of 30 and 100 μm. The results show that the depressed cladding waveguide will be helpful for up-conversion applications in optical integrated devices.

Three possible up-conversion mechanisms for Er$^{3+}$/Yb$^{3+}$ co-doped phosphate glasses are by excited state absorption (ESA), energy transfer (ET) and cross-relaxation (CR). Based on the energy match situation dependency on pumping power and the Er$^{3+}$ and Yb$^{3+}$ ion energy level diagram, the possible up-conversion mechanism was discussed as shown in Figure 6B. Under 980 nm excitation, the pumping laser operating excites electrons from the ground state to the $2F5/2$ excited state through ground state absorption (GSA) of the exciting photons, i.e., $4F7/2$ (Yb$^{3+}$) + a photon $\rightarrow 2F5/2$ (Yb$^{3+}$). After the first excitation to the $2F5/2$ level, a second photon is absorbed by the same ion (ESA), exciting it to $4F7/2$ state, i.e., $2F5/2$ (Yb$^{3+}$) + a photon $\rightarrow$ $4F7/2$ (Er$^{3+}$). For luminescence, there may be two attenuation possibilities: 1) electrons return from the $4F7/2$ state (Yb$^{3+}$) or $4F5/2$ state (Er$^{3+}$) to the ground state by releasing energy; 2) the rapid nonradiative decay from $4F7/2$→$4F9/2$, $4S3/2$ and $2H11/2$ level results in the dominant population at $4F9/2$, $4S3/2$ and $2H11/2$ states. Further electrons at excited levels return to ground level and emit photons.
Compared with the previous phosphate glass optical waveguide, the depressed waveguides prepared by femtosecond laser writing have lower losses [38–40]. The waveguides can support optical transmission in two polarization directions of TE and TM. It can maintain the excellent up-conversion performance of the material in the waveguide core [41]. The waveguide can be designed for a wider range of cross-sectional dimensions.

CONCLUSION

We have successfully fabricated the waveguides of different diameters in the Er, Yb co-doped phosphate glass by direct femtosecond laser writing. The measured modal profiles of waveguides exhibit fundamental transmission at 633 nm wavelengths. Propagation loss as low as 1.8 dB/cm is obtained. The near fundamental mode and multi-mode are supported via controlling the diameter of waveguide. The up-conversion emissions of the waveguides have been obtained under excitation of 980 nm radiations. We observed the emission spectrum and found that the depressed cladding waveguide will be helpful for up-conversion applications in optical integrated devices. These cladding waveguides show good potential for integrated optical circuits and miniature visible fluorescence waveguide devices.

DATA AVAILABILITY STATEMENT

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

AUTHOR CONTRIBUTIONS

YZ performed simulation, sorted out the data and completed the paper writing. CC performed waveguide fabrication and test. LH made the phosphate glass sample. FC supervised the program and gave some suggestions to this paper.

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