Experimental and Theoretical IR Spectra of 2-Nitrosopyridines

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INTRODUCTION

-NITROSO compounds exist in two forms, as monomers and as dimers–azodioxides.[1] Since in the solid state most of them appear as dimers, it is difficult to record and interpret the IR spectra of monomeric forms. This became much easier after the development of method by which the monomers in solid state were prepared by photolysis of azodioxide polycrystals in KBr pellet at cryogenic temperatures (12–23 K).[1,2] Although nitroso monomers have so far routinely been identified by strong IR absorbance in the spectral region near to 1500 cm−1, assigned to N=O stretching vibration,[1] the detailed interpretation of spectra of these compounds has never been reported.

In this work we report FTIR spectra of three 2-nitrosopyridines (1–3, Scheme 1) recorded after the photolysis of their dimers in KBr pellet at 12 K, as well as their identification and partial assignations by the comparison of observed frequencies with the vibrational frequencies calculated by the density functional theory (DFT) approach.

Also, as the DFT results proved to be quite density functional dependent, molecules containing nitroso group emerged as challenging model for testing various functionals’ ability to produce accurate theoretical vibrational frequencies. A number of density functionals and basis sets were employed for the calculation of harmonic and anharmonic vibrational frequencies of three 2-nitrosopyridines (1–3) and nitrosobenzene,[3] and the results compared between themselves and with the experimental frequencies to find the optimal level of theory for the job in hand.
EXPERIMENTAL

Preparation of Compounds
Dimers of 2-nitrosopyridines 1–3 were prepared as reported before.\[4\]

Spectroscopic Methods
For FTIR measurements Bruker Equinox FTIR spectrometer with 1 cm\(^{-1}\) resolution was used. The sample was prepared as standard KBr pellet cooled by Leybold-Heraeus ROK 10–300 cyclic helium cryostat. Photolysis of crystal was conducted by a high-pressure Hg lamp (250 W).

Computational Methods
For geometrical optimizations and calculations of vibrational frequencies three quantum chemical programs were used: Gaussian 16 rev. C.01,\[5\] Orca 4.2.1\[6\] and Q-Chem 5.1.2.\[7\] Most of the calculations were done with Gaussian 16 because it has implemented analytical second derivatives for all its functionals, which significantly speeds up frequency jobs. In four cases harmonic frequencies had to be calculated numerically, SCS-MP2 and B2GP-PLYP-D3(BJ) frequencies with Orca, and \(\omega_{B97X-V}\) and \(\omega_{B97M-V}\) frequencies with Q-Chem.

Geometry optimizations with all three programs were conducted using tighter optimization criteria ('opt=tight' in Gaussian, 'tightopt' in Orca, and equivalent in Q-Chem), and denser integration grids ('int=ultrafine' or 'int=superfine' in Gaussian, 'Grid6' in Orca and 'SG-3' in QChem).

Harmonic vibrational normal modes obtained with Gaussian and Orca were analyzed with program vibAnalysis.\[8\] Based on VMARD method,\[9\] it uses Bayesian linear regression to decompose vibrational normal modes into a small number of the most prominent internal coordinates describing individual atomic motions. Comparison and analysis of harmonic and anharmonic spectra obtained with Gaussian program was also done with GaussView v6.\[10\]

Theoretical infrared spectra were calculated from scaled harmonic or anharmonic frequencies by in-house program using Lorentzian lineshape function with peak half-width at half height 4 cm\(^{-1}\), and plotted with graphing program DataGraph.\[11\]

Scaling factors for obtaining fundamental frequencies from harmonic frequencies were calculated following the procedure described by Martin,\[12\] using frequency dataset from the same paper (modified HFREQ2014) containing 30 molecules with total of 119 experimental fundamental frequencies. Fitting was done with program DataGraph.

RESULTS AND DISCUSSION

Solid-State Photolysis of 2-Nitrosopyridines Dimers

2-NITROSOPYRIDINE (1)
FTIR spectrum of 2-nitrosopyridine dimer recorded in KBr pellet at 12 K exhibits intense signals at 1390 cm\(^{-1}\) and 1409 cm\(^{-1}\) assigned to the symmetric and asymmetric stretching vibrations of the Z-ONNO group\[13\] (Figure 1a). After photolysis by a high-pressure Hg lamp, the intensities of these signals decrease and several new signals appear (Figure 1b). The most prominent of them at 1509 cm\(^{-1}\) can be tentatively assigned to the stretching vibration of N=O group, indicating formation of monomer 1. By warming the sample to the room temperature, the signals assigned to dimer restore their intensity, and the signals of the monomer disappear (Figure 1c).

4-CHLORO-2-NITROSOPYRIDINE (2)
Similarly as with 1, irradiation of dimer of 2 by a high-pressure Hg lamp in KBr pellet at 12 K induces its photodissociation to monomers. FTIR spectra show that starting mixture of Z- and E- dimers also undergoes photodissociation in KBr pellet at 12 K. Intensities of the signals of Z- and E-ONNO groups at 1389 and 1250 cm\(^{-1}\) decrease upon irradiation, whereas the new signal attributed to monomeric N=O group appears at 1507 cm\(^{-1}\). Both Z- and E-dimers are regenerated by thermal dimerization.

Figure 1. FTIR spectra of 1 in KBr pellet (a) at 12 K, (b) at 12 K after photolysis by a high-pressure Hg lamp, and (c) after warming to room temperature.
4-METHYL-2-NITROSYRIDINE (3)

Dimer of compound 3 affords solid-state photo behavior similar to 2-nitrosopyridine and its 4-chloro derivative 2. Intensities of the signals of Z- and E-dimers at 1400 and 1266 cm\(^{-1}\) are reduced upon irradiation and a new intensive band is assigned to stretching vibration of monomeric N=O group appears at 1512 cm\(^{-1}\). Warming the sample to the room temperature regenerates the starting spectrum.

Appearance of monomers in solid phase after cryogenic photolysis opens opportunity to record and examine their yet unknown vibrational spectra. In Figure 2 are given differential FTIR spectra obtained after photolysis of corresponding dimers. All three of them contain a strong band between 1509 and 1512 cm\(^{-1}\), and since it is known that nitroso compounds have N=O bond stretching band in this region,\(^{[1]}\) it is reasonable to assume that irradiation of dimers produces corresponding monomers. The best way to confirm it is to compare theoretical vibrational frequencies of monomers with the frequencies of the bands in experimentally recorded differential spectra.

Calculation of Vibrational Spectra of Monomers

There are two standard ways to calculate vibrational frequencies of molecules. Simpler and more often used is based on harmonic force field approximation, in which potential energy in the vicinity of stationary point is expanded into Taylor series and truncated after second order term. This allows calculation of harmonic vibrational frequencies by diagonalization of the massweighted second order energy derivatives matrix. Due to anharmonicity effects, electron correlation effects, and basis sets incompleteness, frequencies calculated in this way are higher than the experimental fundamental frequencies. The simplest of several ways\(^{[14]}\) to deal with this problem is to multiply calculated harmonic frequencies by a single scaling factor specific for the particular level of theory.

More accurate approach, though rather lengthy one, is to calculate anharmonic corrections to the vibrational energy using second-order vibrational perturbation theory (VPT2).\(^{[15]}\) As it requires third and fourth order energy derivatives with respect to the normal coordinates, the procedure is feasible only if used quantum mechanical model has implemented analytical second derivatives.

We have chosen to calculate vibrational frequencies of monomers using density functional theory. Since preliminary calculations revealed that calculated 2-nitrosopyridines frequencies are quite density functional dependent, a larger number of functionals were tested to find the ones that produce reliable results. Functionals were selected for testing either because of their popularity\(^{[16]}\) and widespread presence in the literature, or because they were well positioned in reliable benchmark studies of density functionals performance.\(^{[17]}\) In the interest of efficiency, for the calculation of harmonic vibrational frequencies we have considered only functionals for which at least analytical first derivatives were available, and for the calculation of anharmonic frequencies functionals with analytical second derivatives as well.

In combination with the selected functionals, several families of basis sets were tested: Pople’s\(^{[18]}\) 6-311+G(d,p) and 6-311+G(2df,2p), Dunning’s cc-pVnZ and aug-cc-pVnZ,\(^{[19]}\) Jensen’s PCSEG-\(n\) and aug-PCSEG-\(n\),\(^{[20]}\) and Karlsruhe def2\(^{[21]}\) basis sets of double, triple and quadruple-\(\zeta\) quality. Comparison of spectra of nitroso compounds calculated with different members of the same basis sets family demonstrated that spectra basically do not change beyond triple-\(\zeta\) level, and do not change with the addition of diffuse functions. Since MP2 calculations were performed as well, the advantage was given to def2-TZVP basis set, whose overall performance was the most consistent of all tested basis sets.

So far, scaling factors for conversion of calculated harmonic to fundamental frequencies for def2 basis sets have been published, to the best of our knowledge, only by Truhlar\(^{[22]}\) and Martin.\(^{[12]}\) Truhlar’s scaling factors for various density functionals and mainly def2-TZVP basis set were based on F38/10 dataset consisting of 38 experimental fundamental frequencies of 15 molecules. Martin’s
scaling factors for MP2, SCS-MP2, and a variety of density functionals were based on modified HFREQ2014 dataset consisting of 119 experimental fundamental frequencies of 30 molecules. The rest of scaling factors used in this paper were determined by ourselves using Martin’s procedure and his dataset given in ref. 12.

### Comparison of Calculated and Experimental Frequencies

The main difference between different functionals are the frequencies and intensities of the normal modes containing N=O stretching vibration. While spectra calculated with the most of double-hybrid functionals look just like the experimental spectra, spectra calculated with other functionals significantly deviate, suggesting that not all functionals are capable of predicting accurate vibrational frequencies of these compounds.

Another potential complication is that 2-nitrosopyridines have two conformers, syn and anti (Scheme 2), making the analysis of their experimental spectra more difficult.

On all tested DFT levels of theory anti conformers of 2-nitrosopyridines 1–3 are for 0.4–0.7 kcal mol⁻¹ more stable than corresponding syn conformers. In Table 1 are given their relative standard Gibbs energies, and barriers for gas-phase interconversion by out-of-plane rotation of N=O group around C–N(=O) bond, all calculated at DSD-PBEP86-D3(BJ)/def2-TZVP level of theory.

Since in solid state pyridine nitrogens in Z-dimer are syn, and in E-dimer anti to each other, one would expect that cryogenic photolysis in both cases produces at least predominantly anti-2-nitrosopyridine (Scheme 3). Especially since at 12 K isomerization is very slow, and the difference in stabilities is sufficient that equilibrium mixture contains virtually no syn conformer.

Scaled harmonic infrared spectra of anti and syn 2-nitrosopyridine calculated at DSD-PBEP86D3(BJ)/def2-TZVP level of theory are very similar. Although there are some differences in frequencies and intensities of main bands, they are not significant enough to be useful for the analysis of the experimental spectra. Anharmonic spectra calculated at the same level of theory are even closer in frequencies, but relative intensities of two conformers are quite different. The fact that the intensities of anti conformer closely resembles intensities of experimental spectrum, and those of syn conformer do not, suggests dominant presence of anti conformer. Additional evidence in favor of the prevailing formation of anti conformer comes from the anharmonic spectra of 4-chloro-2-nitrosopyridine. At DSD-PBEP86-D3(BJ)/def2-TZVP level of theory third strongest band of anti conformer’s spectrum is located at 705 cm⁻¹, close to the experimental frequency of 701 cm⁻¹; while in syn conformer’s spectrum the same band is shifted to 764 cm⁻¹. In experimental spectrum there are no bands around 764 cm⁻¹.

To estimate a degree of similarity between calculated and experimental frequencies, and to minimize a possibility of incorrect assignations, several strongest and

| Relative standard Gibbs energies (in kcal mol⁻¹) at 298.15 K |  |
|------------------------------------------------------------|--|
| 2-nitrosopyridine (R = H)                                   |  |
| anti                                                       | 0.00 | 0.00 |
| syn                                                        | 0.61 | 0.66 |
| TS                                                         | 5.69 | 5.13 |
| 4-methyl-2-nitrosopyridine (R = Cl)                         |  |
| anti                                                       | 0.00 | 0.00 |
| syn                                                        | 0.66 | 0.72 |
| TS                                                         | 5.55 | 4.99 |
| 4-methyl-2-nitrosopyridine (R = CH₃)                        |  |
| Anti                                                       | 0.00 | 0.00 |
| Syn                                                        | 0.39 | 0.62 |
| TS                                                         | 5.85 | 5.14 |
| nitrosobenzene                                             | 0.00 | 7.83 |

![Scheme 2](image)

![Scheme 3](image)
relatively bands in experimental IR spectra of three anti-2-nitrosopyridines (1–3) and previously recorded spectrum of nitrosobenzene[31] were selected for comparison with the calculated frequencies. These eight reference bands are the strong band close to 1500 cm$^{-1}$ in each of four experimental spectra, band at around 1150 cm$^{-1}$ in all spectra except in that of 4-methyl-2-nitrosopyridine, and relatively strong and isolated band at 701 cm$^{-1}$ in the spectrum of 4-chloro-2-nitrosopyridine.

Decision which band in scaled harmonic spectra corresponds to particular reference band was not in all cases straightforward, or even possible. Spectra obtained with several of the density functionals in the N=O stretching region are so different from both the experimental spectra and spectra obtained with other functionals, that any choice would be arbitrary. To help the objective assignments, compositions of normal modes whose scaled frequencies were reasonably close to the frequencies of reference bands were analyzed using the program vibAnalysis,[8,9] and the results compared with the results for other functionals. All of the decisions were thus based upon estimates of normal modes frequencies, compositions and relative intensities.

In calculated spectra of all nitroso compounds, region between 1450 and 1650 cm$^{-1}$ contains several bands, one or two of which are among the strongest bands in spectrum. Vibrational analysis reveals that they belong to normal modes with at least some N=O bond stretching contribution. In most cases, one mode has significantly higher N=O stretching contribution than the others (40–70 %), and often the highest intensity as well. In such cases, these bands were confidently selected for frequency comparison with the frequency of corresponding reference bands. However, in several instances calculated spectrum contains two or even three bands with similar intensities and similar N=O stretching contributions. Although such patterns are in disagreement with the results of other functionals, just for the orientation, their average frequency was nonetheless compared with the reference frequency.

All calculated spectra contain an isolated band in second spectral region, between 1100 and 1150 cm$^{-1}$, which by both frequency and intensity resembles experimental 1150 cm$^{-1}$ band. Vibrational analysis of the corresponding normal mode shows that it is in-plane ring distortion with an addition of 8–14 % of C–N(=O) bond stretching. In case of 4-chloro-2-nitrosopyridine, it also contains 6–8 % of C–Cl bond stretching. Although it is present in the spectra of 4-methyl-2-nitrosopyridine as well, the presence of neighboring bands of similar intensity make assignments uncertain.

And finally, in all calculated spectra of 4-chloro-2-nitrosopyridine there is an easy identifiable, isolated band at around 700 cm$^{-1}$, clearly resembling 701 cm$^{-1}$ band in the experimental spectrum. Vibrational analysis shows that it belongs to in-plane ring deformation normal mode that contains 16–17 % of C–Cl bond stretching.

In Table 2 for each level of theory are given differences between calculated and experimental frequencies for the reference bands, corresponding RMSDs (root mean square deviations), and scaling factors used for conversion of harmonic to fundamental frequencies. Scaling factors are either taken from the literature,[11,22] or are new, determined in this paper. Additional scaling factors, not shown in Table 2, can be found in Supporting information, Table S1.

Actually, harmonic frequencies of all four compounds were calculated with larger number of functionals than shown in Table 2. For instance, in addition to TPSSh-D3(BJ) and DSD-PBE0-P86-D3(BJ), frequencies were also calculated with revTPSSh-D3(BJ) and revDSD-PBE0-P86-D3(BJ). As both revised versions produce marginally larger RMSD values, only the original versions of these functionals are included in Table 2. In almost all examined cases, with the exception of PE0-D3(BJ), dispersion-corrected[23] forms of functionals have slightly lower RMSD values than their uncorrected forms, so the results for B3LYP, TPSSh, B2PLYP, etc., are also not shown in Table 2.

Based on differences between calculated frequencies and frequencies of selected experimental bands, measured by the corresponding RMSDs, density functionals shown in Table 2 can be divided into two groups. The first contains so-called double-hybrid (DH) functionals, belonging to fifth rung of Jacob’s ladder,[24] generally considered as the most accurate and universal methods for calculation of molecular properties.[33] Scaled harmonic frequencies calculated with them, with the exception of B86-DH and PB-E0-DH, agree exceptionally well with the selected experimental frequencies, with average RMSD of just 15 cm$^{-1}$. The smallest RMSD values have DSD-PBE0-P86-D3(BJ) (11 cm$^{-1}$), B2GP-PLYP-D3(BJ) (12 cm$^{-1}$), and mPW2-PLYP-D2 (13 cm$^{-1}$). These excellent results support two conclusions. First, that DH functionals with def2-TZVP basis set are capable of producing accurate experimental frequencies of nitroso compounds, and second, that photolysis of 2-nitrosopyridine dimers indeed produces corresponding monomers.

The second group of functionals, comprising non-DH functionals, have four times higher average RMSD of 65 cm$^{-1}$. The best of them (TPSSh-D3(BJ), B97-D3(BJ) and B3LYP-D3(DJ)) have RMSDs between 32 and 39 cm$^{-1}$, and the rest have average RMSD of 77 cm$^{-1}$. What is the origin of these bad numbers? While average RMSD for 1150 cm$^{-1}$ and 700 cm$^{-1}$ bands for nonDH functionals is probably acceptable (26 cm$^{-1}$), average RMSD for 1500 cm$^{-1}$ bands is 85 cm$^{-1}$, or 107 cm$^{-1}$ without three outliers. As B3LYP-D3(BJ) and B97-D3(BJ) have problems with reproducing the
Table 2. Differences between calculated and experimental reference bands frequencies, corresponding RMSDs, and scaling factors ($\lambda_i$) used for conversion of calculated harmonic to experimental fundamental frequencies.

| Method$^{(a)}$ | $\lambda_i$ | lit.$^{(a)}$ | Compounds and frequencies of selected reference bands (in cm$^{-1}$) | \begin{array}{c|c|c|c|c|c|c}
1 (R = H) & 2 (R = Cl) & 3 (R = Me) & Nitrosobenzene & RMSD \\
1509 & 1150 & 1507 & 1160 & 701 & 1512 & 1503 & 1118 \\
\hline
B3LYP-D3(BJ) & c & T & 43$^{(c)}$ & -34 & 66 & -38 & -22 & 39$^{(c)}$ & 15 & -32 & 39 \\
B97-D3(BJ) & 0.9867 & & 2 & -65$^{(c)}$ & 47 & -48 & -19 & -2 & 0 & -46 & 37 \\
M06-2X & 0.946 & T & 137 & -21 & 142 & -24 & -22 & 92 & 130 & -32 & 91 \\
MN15 & 0.9515 & & 118 & -14 & 122 & -20 & -31 & 115 & 52 & -32 & 77 \\
PBEO & 0.9575 & & 114 & -12 & 116 & -20 & -16 & 111 & 52 & -23 & 73 \\
PWB95-D3(BJ) & 0.9542 & & 103 & -15 & 105 & -23 & -21 & 85$^{(c)}$ & 46 & -26 & 64 \\
TPSSH-D3(BJ) & 0.9686 & & 16 & -32 & 60 & -36 & -24 & 12 & 13 & -32 & 32 \\
$\omega$B97X-D & 0.950 & T & 112 & -18 & 116 & -21 & -18 & 110 & 54 & -24 & 73 \\
$\omega$B97X-V & 0.9539 & & 130 & -9 & 115 & -26 & -18 & 128 & 73 & -18 & 82 \\
$\omega$B97M-V & 0.9578 & & 123 & -11 & 114 & -20 & -17 & 123 & 70 & -21 & 79 \\
\hline
\begin{array}{c}
\text{non-DH functional, scaled harmonic frequencies}
\end{array} \\
\hline
TPSSH-D3(BJ) & & & 19 & -23 & 18 & -23 & -10 & 29 & 28 & -27 & 23 \\
\hline
\begin{array}{c}
\text{MP2/SCS-MP2, scaled harmonic frequencies}
\end{array} \\
MP2 & 0.9575 & M & -68 & -5 & -74 & -19 & -23 & -72 & -104 & -15 & 58 \\
SCS-MP2 & 0.9589 & M & -45 & -6 & -48 & -14 & -22 & -49 & -29 & -14 & 33 \\
\hline
\begin{array}{c}
\text{DH functional, scaled harmonic frequencies}
\end{array} \\
B2GP-PLYL-D3(BJ) & 0.9543 & & -5 & -8 & -6 & -17 & -17 & -10 & -9 & -18 & 12 \\
B2-PLYL-D3(BJ) & 0.9623 & M & -16 & -19$^{(c)}$ & -16 & -24 & -16 & -21 & -48$^{(c)}$ & -23 & 25 \\
DSD-PBEPR6-D3(BJ) & 0.9617 & M & -6 & -7 & -5 & -13 & -16 & -10 & -8 & -17 & 11 \\
mPW2-PLYP-D2 & 0.9572 & 4 & -13 & 4 & -21 & -17 & 1 & -1 & -20 & 13 \\
revDSD-BLYP-D3(BJ) & 0.9565 & & -14 & -16 & -12 & -13 & -16 & -17 & -12 & -16 & 15 \\
PBE-DH & 0.9437 & 109 & -2 & 111 & -9 & -16 & 92$^{(c)}$ & 53 & -18 & 67 \\
PBE-QIDH & 0.9421 & 37 & 0 & 80 & -8 & -17 & 6$^{(c)}$ & 32 & -16 & 41 \\
\hline
\begin{array}{c}
\text{DH functional, anharmonic frequencies}
\end{array} \\
DSD-PBEPR6-D3(BJ) & & 17 & 10 & 18 & 4 & 4 & 18 & 25 & 1 & 15 \\
\hline

$^{(a)}$ Basis set is def2-TZVP.

$^{(b)}$ Scaling factors: T=ref. 20, M=ref. 11, the rest= this paper.

$^{(c)}$ Calculated frequency is an arithmetic mean of two frequencies of similar intensities and, if applicable, of similar

Differences between calculated and experimental reference bands frequencies, corresponding RMSDs, and scaling factors ($\lambda_i$) used for conversion of calculated harmonic to experimental fundamental frequencies. The correct shape of the spectra, and TPSSH-D3(BJ) predicts too high N=O stretching frequency of 4-chloro-2-nitroso-pyridine, it can be concluded that none of non-DH density functionals shown in Table 2 are capable of producing acceptable frequencies of all investigated nitrosoaromatic compounds. To see whether this is just the feature of nitrosoaromatic compounds, we have calculated vibrational frequencies of nonaromatic nitroso compound, CH$_3$NO (Table S2 in Supporting information), for which vibrational spectrum is already known.

Average of RMSDs for N=O stretching band obtained with B3LYP-D3(BJ), B97-D3(BJ) and TPSSH-D3(BJ) functionals is 40 cm$^{-1}$, with other non-DH functionals it is 97 cm$^{-1}$, and with DH functionals (without two outliers) it is just 9 cm$^{-1}$. It can be safely concluded that the origin of the problem is in the N=O bond itself, and has nothing to do with the rest of the molecule. The inability of non-DH (and some DH) density functionals to reproduce the N=O stretching frequency can be ascribed to their inadequate quantum-mechanical description of N=O group, which is probably the

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We have also tested MP2 and SCS-MP2 methods. The RMSD value for MP2 (58 cm⁻¹) is somewhat lower than average non-DH DFT RMSD, but SCS-MP2 (33 cm⁻¹) offers definitive improvement. However, since the result of SCS-MP2 is only in the best non-DH DFT region, there is no real reason to use it.

To see whether there are any significant changes relative to the simpler and faster approach, we have calculated anharmonic spectra of all compounds with the best functional from both groups. From the resulting RMSD values reported in Table 2 it is clear that TPSSH-D3(BJ) anharmonic frequencies represent an improvement relative to the harmonic frequencies, but still not in the range of frequencies obtained with DH functionals. The anharmonic frequencies obtained with DSD-PBE86-D3(BJ) functional produce practically the same RMSD value (15 cm⁻¹) as do the scaled harmonic ones (11 cm⁻¹), leading to the conclusion that tedious calculation of anharmonic frequencies with DH functionals, at least for the compounds investigated in this paper, does not necessarily offer the advantage relative to the simpler approach of multiplying calculated harmonic frequencies with the single scaling factor specific for the level of theory.

There is a feature common to all scaled harmonic frequencies calculated with DH functionals, not counting two flawed frequencies obtained with PBE0-DH and PBE-QIDH. Corresponding differences shown in Table 2 are all negative, with average value of −10 cm⁻¹, meaning that all frequencies are slightly underestimated by one-parameter (1ν) uniform scaling procedure [Eq. (1)]. This suggests that possibly better results can be obtained with linear scaling equation approach \[ \tilde{\nu}_{\text{exp}} = \lambda \tilde{\nu}_{\text{calc}} \] (Eq. (2)).

Parameters \( a \) and \( b \) for DSD-PBE86-D3(BJ)/def2-TZVP level of theory were determined by linear correlation of calculated harmonic and experimental fundamental frequencies of Martin’s dataset, and then used for scaling harmonic frequencies of four molecules examined in this paper. As the new frequency differences are smaller and of both signs, their RMSD value is reduced to 3 cm⁻¹, and their average value of −2 cm⁻¹ is much closer to zero.

There is an excellent agreement between experimental and scaled harmonic spectra, with almost exact match between most of the corresponding bands. Not only in frequencies, but in relative intensities as well. Small number of bands present in the calculated spectra are either missing in the experimental spectra or have smaller intensities than predicted, which in principle might have to do with the interactions in the solid phase. Most notably, these are the bands close to 1600 cm⁻¹, just above strong N=O stretching bands, which belong to skeletal in-plane deformations of aromatic ring, usually with small addition of N=O stretching vibration. Conversely, just two not so strong bands in three experimental spectra do not have apparent counterparts neither in scaled harmonic nor anharmonic spectra – 1131 cm⁻¹ band in the spectrum of 1 and 977 cm⁻¹ band in the spectrum of 3. In other words, scaled harmonic spectra calculated with DSD-PBE86-D3(BJ) DH functional and def2-TZVP basis set can evidently serve as a good starting point for more detailed assignment of experimental spectra of nitrosoaromatic compounds.
CONCLUSIONS

Excellent agreement between experimental and theoretical IR spectra calculated with DH functionals and def2-TZVP basis set, confirms that photolysis of 2-nitrosopyridine dimers at 12 K produces corresponding monomers. The best results were obtained with DSD-PBEP86-D3(BJ), B2GP-PLYP-D3(BJ) and mPW2-PLYP-D2 functionals.

Surprisingly, non-DH functionals, including otherwise very good ωB97X-V[28] and ωB97M-V[29] functionals, have serious problems with normal modes containing N=O bond stretching vibration. As the same holds for non-aromatic nitrosomethane, the origin of the problem is clearly in the nitroso group itself, most likely in its multireference character.

For the compounds investigated in this paper, anharmonic corrections do not necessarily end up in frequencies closer to the experimental. The best scaled harmonic frequencies of 2-nitrosopyridines (1-3) and nitrosobenzene among tested non-DH functionals belong to TPSSH-D3(BJ), B97-D3(BJ) and B3LYP-D3(BJ) functionals and def2-TZVP basis set. While anharmonic corrections improve TPSSH-D3(BJ) frequencies, the effect on frequencies obtained with two other functionals is the opposite one. Similarly, scaled harmonic frequency of nitrosomethylene N=O stretching vibration is closer to the experimental value than its anharmonic equivalent. In the case of the best DH functional, DSD-PBEP86-D3(BJ), and several other DH functionals as well, there are no significant differences between scaled harmonic and anharmonic frequencies and intensities.

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Supplementary Information. Supporting information to the paper is attached to the electronic version of the article at: https://doi.org/10.5562/cca3820.

REFERENCES

[1] H. Vančík, Aromatic C-nitroso Compounds, Springer, Dordrecht, 2013.

[2] H. Vančík, V. Šimunić-Mežniaric, I. Čaleta, E. Meštrović, S. Milovac, J. Phys. Chem. B 2002, 106, 1576–1580. https://doi.org/10.1021/jp0115289

[3] A. Knežević, T. Medančić, S. Milovac, I. Biljan, I. Halasz, H. Vančík, Croat. Chem. Acta 2011, 84, 21–24. https://doi.org/10.5562/cca1714

[4] K. Varga, I. Biljan, V. Tomišić, Z. Mihal, H. Vančík, J. Phys. Chem. A 2018, 122, 2542–2549. https://doi.org/10.1021/acs.jpca.7b12179

[5] Gaussian 16, Revision B.01, M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparrini, F. Egidii, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery, Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Hev, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keigh, R. Kobayashi, J. Normand, R. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman, D. J. Fox, Gaussian, Inc., Wallingford CT, 2016. F. Neeese, F. Wennmohs, U. Becker, C. Riplinger, J. Chem. Phys. 2020, 152, 224108. https://doi.org/10.1063/5.004608

[7] Y. Shao, Z. Gan, E. Epifanovsky, A. T. B. Gilbert, M. Wormit, J. Kussmann, A. W. Lange, A. Behn, J. Deng, X. Feng, D. Ghosh, M. Goldey P. R. Horn, L. D. Jacobson, I. Kaliman, R. Z. Khaliluлин, T. Kús, A. Landau, J. Liu, E. I. Pronyov, Y. M. Rhee, R. M. Richard, M. A. Rohrdanz, R. P. Steele, E. E. Sundstrom, H. L. Woodcock III, P. M. Zimmerman, D. Zuev, B. Albrecht, E. Alguire, B. Austin, G. J. O. Beran, Y. A. Bernard, E. Berquist, K. Brandhorst, K. B. Bravaya, S. T. Brown, D. Casanova, C.-M. Chang, Y. Chen, S. H. Chien, K. D. Crosser, D. L. Crittenden, M. Diederhenhoff, R. A. DiStasio Jr., H. Hop, A. D. Dutoi, R. G. Edgar, S. Fatehi, L. Fusti-Molnar, A. Ghysels, A. Golubeva-Zadorozhnaya, J. Gomes, M. W. D. Hanson-Heine, P. H. P. Harbach, A. W. Hauser, E. G. Hohenstein, Z. C. Holden, T.-C. Jagau, H. J., B. Kaduk, K. Khistyaev, J. Kim, J. Kim, R. A. King, P. Klunzinger, D. Kosenkov, T. Kowalczyk, C. M. Krauter, K. U. Lao, A. Laurent, K. V. Lawler, S. V. Levecchenko, C. Y. Lin, F. Liu, E. Livshits, R. C. Lochan, A. Luenser, P. Manohar, S. F. Manzer, S.-P. Mao, N. Mardirossian, A. V. Marenich, S. A. Maurer, N. J. Mayhall, C. M. Oana, R. Olvares-Amaya, D. P. O’Neill, J. A. Parkhill, T. M. Perrine, R. Pererati, P. A. Pleniz, A. Prociuk, D. R. Rehn, E. Rosta, N. J. Russ, N. Sorguewe, S. M. Sharada, S. Sharma, D. W. Small, A. Sott, T. Stein, D. Stück, Y.-C. Su, A. J. W. Thorn, T. Tsuchimo, L. Vöğt, O. Vydov, T. Wang, M. A. Watson, J. Wenzel, A. White, C. F. Williams, V. Vanovschi, S. Yegehan, S. R. Yost, Z.-Q. You, Y. I. Zhang, X. Zhang, Y. Zhou, B. R. Brooks, G. K. L. Chan, D. M. Chipman, C. J. Cramer, W. A. Goddard III, M. S. Gordon, W. J. Hehre, A. Klann, H. F. Schaefer III, M. W. Schmidt, C. D. Sherrill, D. G. Truhlar, A. Warshell, X. Xia, A. Aspuru-Guzik, R. Baer, A. T. Bell, N. A. Besley, J.-D. Choi, A. Dreuw, B. D. Dunietz, T. R. Furlani, S. R. Gwaltney, C. P. Hsu, Y. Jung.
[8] F. Teixeira, VibAnalysis−Tools for performing vibrational analysis on molecular systems, 2017; https://github.com/teixeirafilipe/vibAnalysis

[9] F. Teixeira, M. N. D. S. Cordeiro, J. Chem. Theory Comput. 2019, 15, 456−470. https://doi.org/10.1021/acs.jctc.8b00439

[10] GaussView, Version 6.0, R. Dennington, T. A. Keith, J. M. Millam, Semichem Inc., Shawnee Mission, KS, 2016.

[11] DataGraph, Version 4.6; https://www.visualdatatools.com/DataGraph

[12] M. K. Kesharwani, B. Brauer, J. M. L. Martin, J. Phys. Chem. A 2015, 119, 1701−1714. https://doi.org/10.1021/jp508422u

[13] B. G. Govenlock, M. J. Maidment, K. G. Orrell, V. Šik, G. M. G. Vasapollo, M. B. Hursthouse, K. M. A. Malik, J. Chem. Soc., Perkin Trans. 2 2000, 2280.

[14] M. A. Palafox, Phys. Sci. Rev. 2018, 3, 1. https://doi.org/10.1515/psr-2017-0184

[15] a) V. J. Barone, Chem. Phys. 2005, 122, 1. https://doi.org/10.1063/1.1824881
b) V. Barone, M. Biczysko, J. Bloino, Phys. Chem. Chem. Phys. 2014, 16, 1759−1787. https://doi.org/10.1039/c4cp04131h

[16] M. Swart, Results for the 2020 edition of the Annual DFT Popularity Poll: Primera División; https://www.marcelswart.eu/dft-poll/res2020_q1.html

[17] L. Goerigk, A. Hansen, C. A. Bauer, S. Ehrlich, A. Najibi and S. Grimme, Phys. Chem. Chem. Phys. 2017, 19, 32184−32215. https://doi.org/10.1039/C7CP04913G

[18] W. J. Hehre, L. Radom, P. v. r. Schleyer, J. A. Pople, Ab Initio Molecular Orbital Theory, New York 1986, John Wiley & Sons.

[19] a) T. H. Jr. Dunning, J. Chem. Phys. 1989, 90, 1007. https://doi.org/10.1063/1.456153
b) A. K. Wilson, D. E. Woon, K. A. Peterson, T. H. Jr. Dunning, J. Chem. Phys. 1999, 110, 7667. https://doi.org/10.1063/1.478678

[20] F. Jensen, J. Chem. Theory Comput. 2014, 10, 1074−1085. https://doi.org/10.1021/ct401026a

[21] F. Weigend, R. Ahlrichs, Phys. Chem. Chem. Phys. 2005, 7, 3297−3305. https://doi.org/10.1039/b508541a

[22] I. M. Alecu, J. Zheng, Y. Zhao, D. G. Truhlar, J. Chem. Theory Comput. 2010, 6, 2872−2887. https://doi.org/10.1021/ct100326h

[23] a) S. Grimme, S. Ehrlich, L. Goerigk, J. Comput. Chem. 2011, 32, 1456−1465. https://doi.org/10.1002/jcc.21759
b) S. Grimme, J. Antony, S. Ehrlich, H. Krieg, J. Chem. Phys. 2010, 132, 154104. https://doi.org/10.1063/1.3382344

[24] J. P. Perdew, K. Schmidt, AIP Conf. Proc. 2001, 577, 1. https://doi.org/10.1063/1.1390175

[25] L. Goerigk, S. Grimme, WIREs Comput. Mol. Sci. 2014, 4, 576−600.

[26] A. J. Barnes, H. E. Hallam, S. Waring, J. R. Armstrong, J. Chem. Soc, Faraday Trans. 2 1976, 72, 1−10. https://doi.org/10.1039/F29767200001

[27] E. K. Dolgov, V. A. Bataev, I. A. Godunov, Int. J. Quant. Chem. 2004, 96, 193−201. https://doi.org/10.1002/qua.10683

[28] N. Mardirossian, M. Head-Gordon, Phys. Chem. Chem. Phys. 2014, 16, 9904−9924. https://doi.org/10.1039/c4cp04374a

[29] N. Mardirossian, M. Head-Gordon, J. Chem. Phys. 2016, 144, 214110. https://doi.org/10.1063/1.4952647