On the Stellar Velocity Distribution in the Solar Neighborhood in Light of Gaia DR2

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Abstract

The aim of this Letter is to contribute to the understanding of the stellar velocity distribution in the solar neighborhood (SN). We propose that the structures on the $U$–$V$ planes, known as the moving groups, can be mainly explained by the spiral arms perturbations. The applied model of the Galactic disk and spiral arms, with the parameters defined by observational data and with pattern speed $\Omega_p = 28.0 \text{ km s}^{-1} \text{kpc}^{-1}$, is the same that allowed us to explain the origin of the Local Arm and the Sun’s orbit trapped inside the corotation resonance (CR). We show that the $U$–$V$ picture of the SN consists of the main component, associated with the CR, and the inner and outer structures, which we could associate with the Hercules and Sirius streams, respectively. The Coma-Berineses and Hyades–Pleiades groups, and the Sun itself, belong to the main part. The substructures of Hercules are formed mainly by the nearby 8/1, 12/1, and even 6/1 inner Lindblad resonances, while Sirius is shaped by the bulk of overlapping outer Lindblad resonances, $\pm 8/1$, $\pm 12/1$, $\mp 16/1$, which are stuck to the CR. This richness in resonances only exists near corotation, which should be of the spiral arms, not of the Galactic bar, whose stable corotation zone is far away from the Sun. The model’s predictions of the velocity distribution match qualitatively and quantitatively the distribution provided by Gaia data release 2.

Key words: galaxies: spiral – Galaxy: kinematics and dynamics – Galaxy: structure – solar neighborhood

1. Introduction

In recent years, great efforts have been dedicated to providing plausible explanations to the moving groups and the rough bimodality that are observed in the velocity distribution of the solar neighborhood (SN; Dehnen 2000; Quillen & Minchev 2005; Antoja et al. 2014). However, these kinematic structures are not yet settled and it seems that, at present, we lack models which could explain the observations as complete as those presented by Gaia data release 2 (DR2; Gaia Collaboration et al. 2018).

The stellar velocity distribution near the Sun shows density clumps, which were first related to the disruption of open clusters (see Antoja et al. 2010, for a review). Nevertheless, the heterogeneity of the ages of these groups (Bensby et al. 2007; Famaey et al. 2008) was in contradiction with this hypothesis, demanding a new scenario. Furthermore, a question that has drawn attention was the apparent bimodality in the $U$–$V$ distribution, in which one of these groups, Hercules, is separated from the main component of the velocity distribution in the SN.

Some dynamical scenarios explain this separation as an imprint of the Galactic bar perturbations (Dehnen 1999, 2000; Bovy 2010; Antoja et al. 2014; Pérez-Villegas et al. 2017). Others study spiral-arm perturbations (De Simone et al. 2004; Quillen & Minchev 2005; Antoja et al. 2011; Quillen et al. 2018a) or a mix of bar+spirals (Antoja et al. 2009). Additionally, transient spiral structures and phase wrapping were recently proposed as an explanation for the $U$–$V$ distribution in the SN (Antoja et al. 2018; Hunt et al. 2018).

Our explanation presented here is based on strong observational evidence, which includes the Sagittarius-Carina and Perseus spiral arms, the Local Arm (Lépine et al. 2017, hereafter Paper I), and the proximity between the spiral corotation and the solar circle (Mishurov & Zenina 1999). We find that the complex structure of the $U$–$V$ plane in the SN can be represented roughly by three components: the main component, associated with the spiral corotation resonance (CR); Hercules, associated with inner Lindblad resonances (ILRs); and Sirius, associated with the bulk of overlapping outer Lindblad resonances (OLRs). Both ILRs and OLRs are high-order resonances, which naturally appear in the corotation neighborhood.

2. Data and Model

The Gaia mission recently released its second set of data (Gaia Collaboration et al. 2018), which provides 6D phase-space coordinates for 7,224,631 stars, i.e., positions in the sky ($\alpha$, $\delta$), parallaxes $\varpi$, proper motions $\mu_\alpha^*$, $\mu_\delta$, and radial velocities $V_r$ (Lindegren et al. 2018), and radial line-of-sight velocities ($V_\perp$) for stars with $G < 13$ mag. We restrict ourselves to stars with parallax errors smaller than 20%. Additionally, we focus on stars inside 1 kpc from the Sun, where using $d = 1/\varpi$ as distance estimate is acceptable. As result, our final catalog contains 3,105,498 stars. In order to convert the positions, parallaxes, proper motions on the sky, and radial velocities of the stars into Cartesian Galactic phase-space positions and velocities, we use the galpy python tools (Bovy 2015).

Our Galactic model is composed of an axisymmetric background derived from the observed rotation curve (Barros et al. 2016), and a four-arm spiral structure, which has Gaussian-shaped azimuthal profiles (Junqueira et al. 2013; Michtchenko et al. 2017). The spiral arms profile, in the rotating frame, is given by

$$\Phi_{\text{sp}}(R, \varphi) = -\zeta_0 R e^{-1/2(1-\cos(m\varphi-f_0(R)))^2/\sigma^2},$$

where $R$ is the Galactocentric distance, and $\varphi$ is the azimuthal angle, measured with respect to the $X$-axis perpendicular to the Sun’s direction. The parameters are: the spiral-arm strength $\zeta_0 = 200.0 \text{ km}^2 \text{s}^{-2} \text{kpc}^{-1}$, the width $\sigma = 10.0 \times \sin(15^\circ)$ kpc...
(i is the pitch angle), and the radial scale-length $\varepsilon_i^{-1} = 4.0 \text{ kpc}$, while $f_{\text{en}}(R)$ is a four-arm logarithmic shape function. More details about the parameters are presented in Paper I. The arm’s width was slightly increased, reducing the chaotic layers of the CR that allows us to clearly visualize the high-order resonances close to the CR.

Our recent study (Paper I) shows that the origin of the Local Arm can be explained by the proximity of the Sun to the CR (Dias & Lépine 2005), provided that the $\Omega_p$-value is in the range $27 \pm 2 \text{ kpc}^{-1}$. We choose here $\Omega_p = 28 \text{ km s}^{-1} \text{kpc}^{-1}$, corresponding to a corotation radius at 8.2 kpc. In our model, the Sun is placed at $R_0 = 8.0 \text{ kpc}$, with a circular velocity of $V(R_0) = 230 \text{ km s}^{-1}$, assuming a solar peculiar motion with respect to the circular orbit of $(U_0, V_0) = (11.1, 12.24) \text{ km s}^{-1}$ (Schönrich et al. 2010).

3. Velocity Distribution in the SN

Using our model, we calculate the dynamical map on the $U-V$ plane, presented in the left panel of Figure 1, for $R = 8.0 \text{ kpc}$ and $\varphi = 90^\circ$. Lighter gray tones on the map show the regular orbits, while increasingly dark tones show the increasing instabilities and chaotic motion. The resonances are recognized as structures surrounded by chaotic layers, associated with resonance separatrices. The purple color highlights the resonant and quasi-resonant domains (or zones of resonance influence). For the calculations, we use integration times of $\sim 12 \text{ Gyr}$, the structures appear clearly on the maps after 2 Gyr ($\sim 10$ orbits around the Galactic center), but we perform longer integrations to get better defined separatrices. The $U-V$ velocity distribution within 150 pc from the Sun provided by Gaia DR2 is shown in right panel of Figure 1.

The initial comparison of the dynamical map (left) with the density distribution (right) in Figure 1 can be done only qualitatively, as the positions of the resonances on the $U-V$ plane depend on $R$. We observe a dominant presence of the CR in the central part of the plane and identify the well-known moving groups, Coma-Berenices and Hyades–Pleiades, as objects belonging to the CR and its zone of influence. The abrupt cutoff on the upper left-side of the main region, observed previously in the data (Antoja et al. 2008), would be explained by the corotation separatrix. The CR separates the $U-V$ plane in the inner and outer regions. The inner region, populated by several high-order ILRs (8/1 being the most influential), is associated with Hercules. Lindblad resonances are also present in the outer region, which we associate with Sirius. The detailed comparison of the $U-V$-structures and Lindblad resonances is presented in Section 5.

4. Velocity Distribution as a Function of $R$

The connection between the corotation/resonances and the moving groups is verified regarding the main features of resonant motion. For instance, our model predicts that the resonant $V$-velocities will depend on Galactocentric distance, as shown in Figure 2, where we present dynamical maps on the $R-V$ planes, for two values of $U$, $-30 \text{ km s}^{-1}$ (top panel) and $+40 \text{ km s}^{-1}$ (bottom panel). On both planes, the resonant $V$-velocities continuously decrease with increasing $R$; consequently, the resonance pattern in the left panel of Figure 1 (calculated for $R = 8.0 \text{ kpc}$) will move up/down for smaller/larger $R$-values. The apparent stable domains of resonant objects on the $U-V$ plane are broader than those shown in Figure 1 left, as the radial distances of the sample are distributed inside 150 pc from the Sun in the right panel of Figure 1.

Absence of objects is expected in the vicinity of the resonance saddle points, where the stable zone is of zero-width and chaotic motion dominates, as shown in Figure 2. The positions of the saddle points and, consequently, of the low-density regions on the $U-V$ planes, change with radius. In addition, we can see that, for $U = -30 \text{ km s}^{-1}$ (top panel), the widths of the stable CR and 8/1 ILR are non-zero for all considered $R$- and $V$-values, excluding only very low $V$-values. In contrast, for $U = +40 \text{ km s}^{-1}$ (bottom panel), the width of the CR is nearly zero over the whole $R-V$ plane, while the
width of the 8/1 ILR drops to zero in the domains $7.4 < R < 8.5$ kpc and $-80 < V < -30$ km s$^{-1}$. The distances between the CR and nearby ILRs increase with increasing Galactocentric distances.

The main features of the CR and Lindblad resonances predicted by our model are compared with the $U$–$V$ stellar velocity distribution in Figure 3, where we show five planes with square bins of 400 pc. The center of each bin, in Cartesian coordinates with respect to the Galactic center, is $(X, Y) = (0, 7.2), (0, 7.6), (0, 8.0), (0, 8.4),$ and $(0, 8.8)$, in kpc. All planes exhibit a dominant presence of the main component associated with effects of the CR, which slowly scrolls down on the $U$–$V$ plane with increasing $R$, as predicted in Figure 2.

5. Hercules’ and Sirius’ Substructures

Gaia DR2 reveals, for the first time, a very complex structure of the inner and outer zones on the $U$–$V$ plane, associated with the Hercules and Sirius streams, respectively. In order to compare the results of our model with the observational data of Figure 3, we present in Figure 4 the dynamical maps on the $U$–$R$ planes, for eight $V$-levels labeled from I to VIII, according to the $V$-values shown by horizontal lines in Figure 3. Each panel is further divided in five $R$-bins, which match the five planes in Figure 3, from (a) to (e), with the corresponding mean $R = \bar{y}$.

The main features that we observe in Figure 4 are the corotation and Lindblad resonances, which shift continuously toward smaller $R$ for increasing $V$ (a behavior already observed in Figure 2). Consequently, the CR dominates the region between the IV- and V-levels on the maps in Figure 4. Each resonance has a specific topology of its domain around stable center and saddle points, which is asymmetric with respect to the $U$-component. Assuming that these dynamical features give rise to structures in the stellar velocity distribution, we look for

Figure 2. Dynamical maps of the CR and nearby resonances as functions of $R$, on the $R$–$V$ planes ($\varphi = 90^\circ$), calculated for $U = -30$ km s$^{-1}$ (top panel) and $U = +40$ km s$^{-1}$ (bottom panel).

Figure 3. Density portraits of the $U$–$V$-plane, as functions of Galactocentric distance; five consecutive zones are labeled from (a) to (e). Each panel shows the velocity distribution in a square of 400 pc on each side. The center of each bin in Cartesian coordinates is indicated in the bottom-left corner, with $X = 0$ for all bins. The color scale is the same as in the right panel of Figure 1. Constant-$V$ dashed lines are denoted as I(−66), II(−50), III(−40), IV(−27), V(0), VI(+10), VII(+20), and VIII(+30) km s$^{-1}$.
their identification in Figure 3. We expect to observe the maximal density of the resonant population in the resonance domain, and the minimal density in the proximity of the saddle points.

5.1. Hercules

Because the Hercules stream is an inner structure, we select four negative \( V \)-levels below the CR (I to IV in Figure 3) to characterize this group. All corresponding panels in the top row of Figure 4 show the presence of the ILRs, from which 8/1 is strongest. Far from the Sun, at \( R = 7.2 \) kpc, the 8/1 ILR already appears at the IV-level. Accompanied by the weaker 6/1 ILR, it produces an \( U-V \)-substructure between the III- and IV-levels, visible in Figure 3(a). Approaching the Sun at \( R = 7.6 \) kpc, this substructure becomes very dense and splits in two horizontal branches, stronger at III-level and weaker at IV-level, clearly observed in Figure 3(b). The corresponding radial bin, (b) in Figure 4, associates the stronger branch to the 8/1 ILR, with its \( U \)-extension ranging approximately from \(-70 \) km s\(^{-1}\) to 30 km s\(^{-1}\), as limited by the proximity of the 8/1 ILR saddle points. The weaker branch is associated, in panel IV(b) of the top row of Figure 4, with the 12/1 ILR (with a remote contribution from the 8/1 ILR).

In the SN, bin (c) in Figure 3, the two branches of Hercules scroll down, being associated now with levels II and III. The density of both decreases as the resonant domains of the corresponding 8/1 and 12/1 ILRs diminish (panels II(c) and III(c) in Figure 4 top). The remnants of these branches are seen in Figure 3(d), which are represented by lines I and II, respectively. For larger radii, the density of Hercules is very small, although we can still observe the tenuous group around \( V = -70 \) km s\(^{-1}\) in panel (e) of Figure 3. The resonant pattern...
of Hercules is therefore consistent with the 8/1 and 12/1 ILRs, being the stronger and weaker branches, respectively.

5.2. *Sirius*

The Sirius stream is an outer structure with respect to the CR, and we select four positive $V$-levels ($V$ to VIII in Figure 3) to characterize this group. The corresponding panels in the bottom row of Figure 4 show the presence of the OLRs, from $-8/1$ to $-16/1$, and higher. Contrasting with high-order ILRs, these resonances are dense and appear to be stuck to the CR, mainly in the V- and VI-levels in bottom row of Figure 4. This peculiar distribution of the OLRs may produce a dynamical phenomenon known as overlapping resonances, which commonly occurs in the vicinity of CRs. The overlap produces an overdensity of objects that we observe in the $U$-$V$-distribution just above the CR in Figure 3. Indeed, Sirius already appears strongly in panel (b) between V- and VII-levels, has maximal density in panels (c) and (d) around the V-level and scrolls down following the CR in (e). The analysis of the corresponding panels in the bottom row of Figure 4 shows that Sirius, located very close to the large separatrix of the CR, is strongly influenced by these resonances in such a way that, for large radii, the separation between them is no longer identifiable (e.g., bins (d)–(e) in Figure 3). Moreover, the influence of the CR may explain the asymmetry with respect to the $U$-velocity contour of Sirius. Perturbations induced by the central bar could also break the symmetry of the $U$-distribution. Precise numerical simulations of the $U$–$V$ stellar distribution could shed light on this issue.

5.3. *Other Structures*

A tenuous substructure below Hercules is observable in panels (a)–(c) of Figure 3. According to panels I(b), II(a), and III(a) in the top row of Figure 4, this group may be associated with the 6/1 ILR. The weak output of the 6/1 ILR is due to the fact that its order is not a multiple of 4, but it is still amazing to know that stars evolving inside this resonance visit the SN from the inner Galactic region (3–4 kpc).

The strong 4/1 OLR appears in panels VI(e), VII(d, e), and VIII(c, d) of the bottom row of Figure 4. We see a small excess of stars in the corresponding upper parts of the bottom panels of Figure 3. Next, we explain the relatively low population of this resonance due to the nature of their orbits.

5.4. *Resonant Orbits*

To complete the analysis, we construct the surfaces of section (SoS) along one energy level, which covers almost all of the possible regimes of motion in the SN. Figure 5 shows the SoS and four representative orbits of stars that visit the SN. The SoS plot shows an eight-island chain (red) below the CR and the corresponding 8/1 ILR orbit in the bottom-left panel. In this case, stars spend most of their time in the SN, enhancing the density in this region. We highlight the 12-island chain of the 12/1 ILR by cyan color and show its orbit in the bottom-right plot.

In the outer region, connected with Sirius, the 8/1 OLR (red chain) appears above the CR in Figure 5, accompanied by the 12/1 and 16/1 OLRs; its orbit is shown on the top-left panel. The four-island cyan chain and its orbit on the top-right panel show that the 4/1 OLR objects come from the Galaxy’s confines, about 13–14 kpc. The 4/1 OLR is very strong, however, we do not observe the related high-density streams on the maps in Figures 3(d), (e). In part, this is due to the fact that these objects spend most of their time far away from the Sun, where their velocities are smaller. To obtain the metallicity of these stars, we select 531 objects from the Rave-TGAS surveys (Lindegren et al. 2016; Kunder et al. 2017) in the interval $30 < V < 100 \, \text{km} \, \text{s}^{-1}$ and find that $\sim40\%$ of this sample evolve inside the 4/1 OLR. From these, the majority ($\sim85\%$) consists of metal-poor stars, with metallicity being $\sim30\%$ in the range $-1.35 < [\text{Fe/H}] < -0.5$ and $\sim55\%$ in the range $-0.5 < [\text{Fe/H}] < 0$. This result reinforces the theory that metal-poor stars come from the outer parts of the Galaxy (Hattori et al. 2018; Quillen et al. 2018b).

6. *Summary*

We presented a novel approach to explain the distribution of stars on the $U$–$V$ plane of the SN, based on a spiral arms
dynamical model. Our basic assumptions are: (i) the proximity between the spiral corotation radius and the solar circle, and (ii) the dynamical stability of the Local Arm and the Sun. The suggestion that both the Local Arm and the Sun evolve inside the spiral CR yields natural constraints on the magnitude of the Galaxy’s spiral pattern speed. For the adopted rotation curve and solar position (Paper I), we chose $\Omega_0 = 28$ km s$^{-1}$ kpc$^{-1}$, which results in a corotation radius of 8.2 kpc.

The comparison of the stellar velocity distribution provided by Gaia DR2 with the dynamical maps of the $U$–$V$ plane shows that it is related to the CR and nearby Lindblad resonances. It is important to state that, in general, resonances modulate qualitatively the dynamics in their environment: they capture and trap stars inside of the stable resonant zones, enhancing the density, and deplete regions close to saddle points and separatrices. Corotation is special among the resonances, being stronger and having a wider zone of influence surrounded by many high-order resonances.

That is exactly what we observe on the $U$–$V$-plane (right panel) in Figure 1: the central region of enhanced density, known as main component, is the CR stable zone according to the dynamical map of the same plane (left panel). The notable moving groups, Coma-Berenices and Hyades–Pleiades, are inside of the CR, together with the Sun, located at the origin of the plane. Immediately above the main component, for $V > 0$, we can observe the Sirius group, whose origin is related to the overlapping high-order OLRs stuck to the CR. Finally, the strong 4/1 OLR is responsible for distant stars to visit the SN.

In the region below the main component, at $V < 0$, the nearby resonances are separated sufficiently on the $U$–$V$ plane to create several substructures well defined by Gaia DR2. The most prominent is the Hercules group related mainly to the 8/1 and 12/1 ILRs. The overdensities produced by the inner resonances are explained by the fact that the objects, coming from the inner Galaxy, spend longer times at the outer edges of their orbits in the SN, because their velocities are smaller there. This kinematics produces the effect of a bimodal distribution on the $U$–$V$ plane.

Despite the simplicity of our model, we are able to explain quantitatively the SN $U$–$V$ structure and its changes with different Galactic radii. The results presented here can be considered as an additional test, which confirms the robustness of our model. They are also another strong indicator that the spiral corotation radius lies near the solar circle, as it is only in the neighborhood of the CR that this abundance of Lindblad resonances can be observed. The bar’s CR has a minor role because, due to its spatial orientation, the $L_4$-points of the bar are distant from the Sun, as shown in Michtchenko et al. (2018).

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