We study the relation between the Laplacian associated to an odd metric on a supermanifold and harmonic superfunctions, through the application of the calculus of variations to a supersymmetric sigma model.

Keywords: Berezinian sheaf, harmonic superfunctions, supermanifolds, supersigma models.

1. Introduction

In the last years (since the apparition of the Batalin-Vilkovisky quantization method) there has been a certain amount of interest on Laplace operators of odd parity on supermanifolds, mainly in connection with their use in the quantization of non Abelian gauge theories (BRST symmetries, see [19,16,12].) These operators have been related to odd Poisson structures and odd divergences, modular classes, etc. (see [20,9,8].) But we want to focus here on the “Riemannian geometric” side of the problem, and the relation between Laplace’s equation and harmonic functions (for a physical interpretation and relation to string theory, see [14].) In the classical setting, harmonic functions are obtained as follows. Consider two fixed Riemannian manifolds \((M, g)\) and \((N, h)\). Given a mapping \(\phi: (M, g) \to (N, h)\), we can take the pullback of the metric \(h\) by \(\phi\) and then raise an index with the inverse metric \(g^{-1}\). If \(A = g^{-1} \cdot \phi^*h\) is the resulting bundle endomorphism \(A: TM \to TM\), its trace \(trA\) is a function on \(M\) which can be integrated (with respect to the Riemannian volume of \(g\)) to give the functional

\[
\phi \mapsto \int_M (trA) \ rvol_g. 
\]
Critical points of this functional are called harmonic mappings (see [5,4,15].) In the particular case \( N = \mathbb{R} \) they are also called harmonic functions, and the Euler-Lagrange equations for the above functional are precisely

\[ \triangle \phi = 0, \]

so that harmonic functions are the solutions to Laplace equations (where the Laplacian is understood as \( \text{div} \circ \text{grad} \).) In Physics, the construction we have described is known as a sigma model, and harmonicity plays a very important rôle in their study. In this paper, we want to show that the relation between harmonicity and the Laplacian is preserved when passing to the supermanifold setting, and this is done by applying the methods of the calculus of variations to the sections of a graded submersion representing supersigma models.

We would like to stress two features of our approach:

1. We use odd metrics.
2. The superlaplacian \( \triangle \) is not trivial, due to the fact that we use the Berezinian sheaf to define it (compare this with the geometric construction in [10], where any superfunction is harmonic).

To make the paper relatively self-contained, we include two brief sections on the Berezinian module and graded metrics on supermanifolds. Then we write our Lagrangian for the supersymmetric sigma model and associate to it a variational problem, so we can apply the techniques of the graded calculus of variations (for a detailed account of these techniques, and general definitions, notations and results on supermanifold theory, see [12,13] and references therein).

2. \( \mathbb{Z}_2 \)-graded metrics

Let \((M, \mathcal{A}) = (M, \wedge \mathcal{E})\) be an \((m|n)\)-dimensional supermanifold given in Batchelor’s form, so \( \mathcal{E} = \Gamma(E) \) where \( E \to M \) is a vector bundle. We denote by \( \{ \partial / \partial x^k \}_{k=1}^n \) a local frame on \( M \), and by \( \{ x^{-j} \}_{j=1}^n \) a local frame of sections of \( E \). By \( \sim \) we will understand the structural morphism of the supermanifold, \( \sim : \mathcal{A} \to \mathcal{C}^\infty(M) \).

As in [11], a \( \mathbb{Z}_2 \)-graded metric (or supermetric) on \( \wedge \mathcal{E} \) is understood to be a graded symmetric, non degenerate \( \wedge \mathcal{E} \)-bilinear map \( G : \text{Der} \wedge \mathcal{E} \times \text{Der} \wedge \mathcal{E} \to \wedge \mathcal{E} \), whose action on a pair \((D_1, D_2) \in \text{Der} \wedge \mathcal{E} \times \text{Der} \wedge \mathcal{E}\) is denoted \( \langle D_1, D_2; G \rangle \) (for another approach to supermetrics, based on principal superbundles, see [13].) We have

1. \( \langle \alpha D_1, D_2; G \rangle = \alpha \langle D_1, D_2; G \rangle, \quad \forall \alpha \in \wedge \mathcal{E} \).
2. \( \langle D_1, D_2; G \rangle = (-1)^{|D_1||D_2|} \langle D_2, D_1; G \rangle \).
3. The map \( G^\circ : D \to \langle D; , ; G \rangle \) is an isomorphism between the \( \wedge \mathcal{E} \)-modules \( \text{Der} \wedge \mathcal{E} \) and \( \text{Hom}(\text{Der} \wedge \mathcal{E}, \wedge \mathcal{E}) = \Omega^1_G(M) \). The inverse of this isomorphism is denoted by \( G^\circ \), as in the non-graded setting. Also, the induced metric on \( \Omega^1_G(M) \) is denoted by \( G^{-1} \).
Note that these conditions imply \( \langle D_1, \alpha D_2; G \rangle = (-1)^{|\alpha||D_1|} \langle D_1, D_2; G \rangle \). Also, we remark that graded forms have a right \( \wedge \mathcal{E} \)-module structure, so we will take care in writing the \( \wedge \mathcal{E} \) factors to the left of graded vector fields and to the right of graded forms.

A graded metric is even (resp. odd) if, \( |\langle D_1, D_2; G \rangle| + |D_1| + |D_2| \equiv 0 \text{mod} 2 \) (resp. \( \equiv 1 \text{mod} 2 \)). In both cases the graded metric is called homogeneous. The associated concept of a graded connection is defined analogously to the non-graded case.

We focus our attention to graded metrics of second-order depth. This means that there is a graded basis (a set of generators of the locally finite \( \wedge \mathcal{E} \)-module \( \text{Der} \wedge \mathcal{E} \)) \( \{ D_\alpha \}_{\alpha=1}^{m+n} \) such that, \( \langle D_i, D_j; G \rangle \in \sum_{0 \leq k \leq 2} \wedge^k \mathcal{E} \). In \([11, \text{Proposition 4.1}]\) the following is proved:

Let \( \langle \cdot, \cdot \rangle \) be a homogeneous graded metric of second order depth which is adapted to the canonical splitting of \( \wedge \mathcal{E} \). A connection \( \nabla \) exists on \( M \) such that, \( \langle \cdot, \cdot; G \rangle = \begin{cases} \begin{pmatrix} g & 0 \\ 0 & \omega \end{pmatrix} & \text{with } \nabla g = 0 \text{ (even case)} \\ \begin{pmatrix} 0 & \kappa \\ \kappa^t & 0 \end{pmatrix} & \text{(odd case,)} \end{cases} \) with respect to the basis \( \{ \nabla \frac{\partial}{\partial x^k}, i_{\frac{\partial}{\partial x^k}} \} \) of \( \text{Der} \wedge \mathcal{E} \), where \( g \) is a metric on \( TM \), \( \omega \) is a symplectic form on \( M \), and \( \kappa: \mathcal{X}(M) \to \mathcal{E}^* \) is a non-degenerate linear map.

If \( d^G \) denotes the graded exterior derivative acting on the algebra of graded forms \( \Omega_G(M) \)—in order to distinguish it from the usual exterior derivative \( d \), which acts on \( \Omega(M) \)—and \( \{ d^G x^k, d^G x^{-j} \} \) is the dual basis on \( \Omega^1_G(M) \) to the basis of derivations \( \{ \nabla \frac{\partial}{\partial x^k}, i_{\frac{\partial}{\partial x^k}} \} \), we can write

\( \langle \cdot, \cdot; G \rangle = d^G x^i \otimes d^G x^{-j} \cdot \kappa_{ij} + d^G x^{-i} \otimes d^G x^j \cdot \kappa_{ji} \).

for any odd metric on \( (M, \wedge \mathcal{E}) \). Note that, in the case \( \mathcal{E} = \Gamma(TM) \), we can make the identifications \( \nabla \frac{\partial}{\partial x^k} = \frac{\partial}{\partial x^k} \) and \( i_{\frac{\partial}{\partial x^k}} = \frac{\partial}{\partial x^k} \) as differential operators on the \( C^\infty(M) \)-algebra \( \wedge \mathcal{E} \).

**Example 1.** Consider the linear supermanifold \( \mathbb{R}^{1|1} = (\mathbb{R}, \Omega(\mathbb{R})) \) with global supercoordinates \( \{ t, \tau \}, |t| = 0, |\tau| = 1 \). The canonical odd Euclidean supermetric can be written as

\( Q = d^G t \otimes d^G \tau + d^G \tau \otimes d^G t \).

Note that \( \mathbb{R}^{1|1} \) admits no even supermetrics (this would require the base manifold being even dimensional).
Let $\mathcal{E} = \Gamma(TM)$ and let $g$ be a metric on $M$. A second-order depth odd supermetric $G$ on $(M, \wedge \mathcal{E})$ can naturally be defined by simply taking as $\kappa$ the induced isomorphism $g : \mathcal{X}(M) \rightarrow \Omega^1(M)$. In local coordinates:

$$G = d^G x^i \otimes d^G x^{-j} \cdot g_{ij} + d^G x^{-i} \otimes d^G x^j \cdot g_{ji},$$

(3)

where $g_{ij}$ is the matrix of $g$ with respect to the local frame $\left( \frac{\partial}{\partial x^i} \right)_{k=1}^m$ on $M$.

### 3. The Berezinian sheaf and divergence

Let $(M, \mathcal{A})$ be a graded manifold, of dimension $(m|n)$, and let $P^k(\mathcal{A})$ be the sheaf of graded $k$–order differential operators of $\mathcal{A}$. This is the submodule of $\text{End}(\mathcal{A})$ whose elements $P$ verify

$$[[P, a_0], a_1], ..., a_k] = 0,$$

for all $a_0, ..., a_k \in \mathcal{A}$ (here we identify an $a \in \mathcal{A}$ with the endomorphism $b \mapsto ab$).

One has that if $\{x^i, x^{-j}\}_{1 \leq i \leq m}$ are supercoordinates for a splitting neighborhood $U \subset M$, $P^k(\mathcal{A}(U))$ is a free module (for both structures, left and right) with basis

$$\frac{\partial^{[\alpha]}}{\partial x^{\alpha}} \circ \frac{\partial^{[\beta]}}{\partial x^{\beta}} = \left( \frac{\partial}{\partial x^i} \right)^{\alpha} \circ \left( \frac{\partial}{\partial x^m} \right)^{\alpha} \circ \left( \frac{\partial}{\partial x^{-1}} \right)^{\beta} \circ \cdots \circ \left( \frac{\partial}{\partial x^{-n}} \right)^{\beta},$$

where $|\alpha| + |\beta| \leq k$.

Let us consider the sheaf $P^k(\mathcal{A}, \Omega^m_G) = \Omega^m_G \otimes_{\mathcal{A}} P^k(\mathcal{A})$, of $m$–form valued $k$–th order differential operators of $\mathcal{A}$. For every open subset $U \subset M$, let $K^n(U)$ be the set of operators $P \in P^n(\mathcal{A}(U), \Omega^m_G(U))$ such that for every $a \in \mathcal{A}(U)$ with compact support, there exists an ordinary $(m – 1)$–form of compact support, $\omega$, fulfilling $P(a) = d\omega$. We observe that $K^n$ is a submodule of $P^n(\mathcal{A}, \Omega^m_G)$ for its right structure, so we can take quotients and obtain the following description of the Berezinian sheaf:

$$\text{Ber}(\mathcal{A}) = P^n(\mathcal{A}, \Omega^m_G) / K^n.$$ (1)

According to this description, a local basis of $\text{Ber}(\mathcal{A})$ can be given explicitly: If $\{x^i, x^{-j}\}_{1 \leq j \leq m}$ are supercoordinates for a splitting neighborhood $U \subset M$, the local sections of the Berezinian sheaf are written in the form

$$\Gamma_U(\text{Ber}(\mathcal{A})) = \left[ d^G x^1 \wedge \cdots \wedge d^G x^m \otimes \frac{\partial}{\partial x^{-1}} \circ \cdots \circ \frac{\partial}{\partial x^{-n}} \right] \cdot \mathcal{A}(U),$$ (2)

where $[\ ]$ stands for the equivalence class modulo $K^n$.

Now, if $X$ is a graded vector field, it is possible to define the notion of graded Lie derivative of sections of the Berezinian sheaf with respect to $X$. This is the mapping

$$\mathcal{L}^G_X : \Gamma(\text{Ber}(\mathcal{A})) \rightarrow \Gamma(\text{Ber}(\mathcal{A}))$$ (3)

given by

$$\mathcal{L}^G_X [\eta^G \otimes P] = (-1)^{|X||\eta^G \otimes P| + 1} [\eta^G \otimes P \circ X],$$ (4)
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for $X$ and $g^G \otimes P$ homogeneous.

This Lie derivative, has the properties that one would expect (cfr. the treatment in [3], Vol. 1 pg. 83):

1. For homogeneous $X \in \text{Der}(A)$, $\xi \in \Gamma(\text{Ber}(A))$ and $a \in A$,
   \[ \mathcal{L}_X^G(\xi \cdot a) = \mathcal{L}_X^G(\xi) \cdot a + (-1)^{|\xi||\xi|} \xi \cdot X(a). \]

2. For homogeneous $X \in \text{Der}(A)$, $\xi \in \Gamma(\text{Ber}(A))$ and $a \in A$,
   \[ \mathcal{L}_a^G \cdot X(\xi) = (-1)^{|X||\xi|} \mathcal{L}_X^G(\xi \cdot a). \]

3. If $\xi_{x^i,x^{-j}} = [d^G x^1 \wedge \cdots \wedge d^G x^m \otimes \frac{\partial}{\partial x^i} \circ \cdots \circ \frac{\partial}{\partial x^m}]$ is the local generator of the Berezinian sheaf on a system of supercoordinates $\{x^i, x^{-j}\}_{1 \leq i \leq m}$, then
   \[ \mathcal{L}_a^G \frac{\partial}{\partial x^i}(\xi_{x^i,x^{-j}}) = \mathcal{L}_a^G \frac{\partial}{\partial x^{-j}}(\xi_{x^i,x^{-j}}). \]

We can now introduce the notion of Berezinian divergence: Let $(M,A)$ be a graded manifold whose Berezinian sheaf is generated by a section $\xi$. The graded function $\text{div}^B(X)$ given by the formula (for homogeneous $X$)
\[ \mathcal{L}_X^G(\xi) = (-1)^{|X||\xi|} \xi \cdot \text{div}^B(X) \]
is called the Berezinian divergence of $X$ with respect to $\xi$. When there is no risk of confusion, we will write simply $\text{div}^B(X)$.

4. (1|1)-supersymmetric sigma model

Below, we consider a model with target $\mathbb{R}^{1|1}$, that is, a mapping (scalar superfield) $\sigma: (M, \Lambda^E) \to \mathbb{R}^{1|1}$, where we consider an odd supermetric on $(M, \Lambda^E)$ as in Example 2 and the canonical metric of Example 1 on $\mathbb{R}^{1|1}$.

This imply that $E \cong \Gamma(TM)$, but this is not a great loss of generality: If $(M, \Lambda^E)$ admits an odd metric, the existence of the non-degenerate pairing $\kappa: X(M) \to E^*$ implies that the dimension of the supermanifold is $(n|n)$.

The mapping $\sigma$ can be viewed as a section of the graded submersion
\[ p: \mathbb{R}^{1|1} \times (M, \Lambda^E) \to (M, \Lambda^E), \]
to which we associate a super-Lagrangian $L \in A_{J^1_G(p)}$ (see [3] for the details of this construction) proceeding by analogy with the non-graded case (which is that of harmonic functions). The coordinates in the bundle of superjets $J^1_G(p)$ will be denoted $\{t, \tau, x^i, x^{-j}, t_i, t_{-j}, \tau_i, \tau_{-j}\}$. Thus, by taking the pull-back (we write $x^\alpha$ collectively for $x^i$ and $x^{-j}$, i.e., $\alpha$ runs from $-n,...,-1$ to $1,...,m$) of (2) along $\sigma$, we obtain
\[ \sigma^* Q = (-1)^{\alpha\beta} d^G x^\alpha \otimes d^G x^\beta \cdot \frac{\partial (t \circ \sigma)}{\partial x^\alpha} \frac{\partial (\tau \circ \sigma)}{\partial x^\beta} + (-1)^{\gamma+1} d^G x^\gamma \otimes d^G x^\delta \cdot \frac{\partial (\tau \circ \sigma)}{\partial x^\gamma} \frac{\partial (t \circ \sigma)}{\partial x^\delta}. \]
As $| - j | = 1, | k | = 0$, and so on, we could write matricially:

$$\sigma^* Q = \left( \begin{array}{ccc} \frac{\partial (t \circ \sigma)}{\partial x^j} & \frac{\partial (t \circ \sigma)}{\partial x^k} & \frac{\partial (t \circ \sigma)}{\partial \sigma} \\ \frac{\partial (\tau \circ \sigma)}{\partial x^j} & \frac{\partial (\tau \circ \sigma)}{\partial x^k} & \frac{\partial (\tau \circ \sigma)}{\partial \sigma} \\ \frac{\partial (\tau \circ \sigma)}{\partial x^j} & \frac{\partial (\tau \circ \sigma)}{\partial x^k} & - \frac{\partial (\tau \circ \sigma)}{\partial \sigma} \end{array} \right).$$

Also, from (3) we have $G^{-1} = g^{ij} : \frac{\partial}{\partial x^j} \otimes \frac{\partial}{\partial x^i} + g^{kl} : \frac{\partial}{\partial x^k} \otimes \frac{\partial}{\partial x^l}$, or equivalently, $G^{-1} = \left( \begin{array}{cc} 0 & g^{ij} \\ g^{ij} & 0 \end{array} \right)$.

We would like to evaluate the action of $G^{-1}$ on $\sigma^* Q$ in such a way that the functorial correspondence with the composition of graded morphisms be preserved. To this end, we must use either the action of the supertrace of this supermatrix, we would obtain 0, as the result is the following:

$$G^{-1} \cdot \sigma^* Q = \left( \begin{array}{ccc} g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^j} + \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right) & \frac{\partial (t \circ \sigma)}{\partial \sigma} \\ g^{ij} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^j} + \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right) & \frac{\partial (\tau \circ \sigma)}{\partial \sigma} \\ \frac{\partial (t \circ \sigma)}{\partial x^j} & \frac{\partial (t \circ \sigma)}{\partial x^k} \end{array} \right).$$

If we were to take the supertrace of this supermatrix, we would obtain 0, as the following computation shows:

$$\text{Str}(G^{-1} \cdot \sigma^* Q) = \text{Tr} \left( g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^j} + \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right) \right)$$

$$- \text{Tr} \left( g^{ij} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^j} + \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right) \right)$$

$$= g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^j} \right) + g^{ij} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^j} \right)$$

$$- g^{kl} \left( \frac{\partial (t \circ \sigma)}{\partial x^k} \right) + g^{kl} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right)$$

$$= g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^j} \right) + g^{ij} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^j} \right)$$

$$+ g^{kl} \left( \frac{\partial (t \circ \sigma)}{\partial x^k} \right) - g^{kl} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right)$$

$$= g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^j} \right) + g^{ij} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^j} \right)$$

$$- g^{kl} \left( \frac{\partial (t \circ \sigma)}{\partial x^k} \right) + g^{kl} \left( \frac{\partial (\tau \circ \sigma)}{\partial x^k} \right)$$

$$= g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^j} \right) - g^{kl} \left( \frac{\partial (t \circ \sigma)}{\partial x^k} \right)$$

Due to this, instead of $\text{Str}$ we will take the usual contraction of $(1, 1)$-forms.
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supertensors, which amounts to
\[
C^1_{1}(G^{-1} \cdot \sigma^*Q) = Tr \left( g^{ij} \left( \frac{\partial}{\partial x^i} \frac{\partial (t \circ \sigma)}{\partial x^j} + \frac{\partial (\tau \circ \sigma)}{\partial x^i} \frac{\partial (t \circ \sigma)}{\partial x^j} \right) \right) + \sum_{k} g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^i} \frac{\partial (t \circ \sigma)}{\partial x^j} \right) \frac{\partial (\tau \circ \sigma)}{\partial x^k} \frac{\partial (t \circ \sigma)}{\partial x^k} + \frac{\partial (\tau \circ \sigma)}{\partial x^i} \frac{\partial (t \circ \sigma)}{\partial x^j} \frac{\partial (t \circ \sigma)}{\partial x^i} \frac{\partial (\tau \circ \sigma)}{\partial x^j} \right) \frac{\partial (t \circ \sigma)}{\partial x^k} \frac{\partial (t \circ \sigma)}{\partial x^k}.
\]

That is, we have arrived at the result
\[
\frac{1}{2} C^1_{1}(G^{-1} \cdot \sigma^*Q) = g^{ij} \left( \frac{\partial (t \circ \sigma)}{\partial x^i} \frac{\partial (t \circ \sigma)}{\partial x^j} + \frac{\partial (\tau \circ \sigma)}{\partial x^i} \frac{\partial (\tau \circ \sigma)}{\partial x^j} \right),
\]
and this suggests to take the following superlagrangian \( L \in A_{JI}(\rho) \) to study the supersymmetric sigma model (where we introduce the obvious notations \( t_{\alpha} = \frac{\partial (t \circ \sigma)}{\partial x^\alpha} \), \( \tau_{\beta} = \frac{\partial (\tau \circ \sigma)}{\partial x^\beta} \)):

\[
L = g^{ij}(t_{\alpha} \tau_{\beta} + t_{\beta} \tau_{\alpha}).
\]

**Remark 3.** The superlagrangian \( L \) is homogeneous of even degree. In Physics literature, it is common to write \( t \circ \sigma = \phi \) and \( \tau \circ \sigma = \psi \), so equation (1) would be

\[
\frac{1}{2} C^1_{1}(G^{-1} \cdot \sigma^*Q) = g^{ij} \left( \frac{\partial \phi}{\partial x^i} \frac{\partial \psi}{\partial x^j} + \frac{\partial \phi}{\partial x^j} \frac{\partial \psi}{\partial x^i} \right)
\]

and the Lagrangian density

\[
L = g^{ij}(\phi_{\alpha} \psi_{\beta} + \phi_{\beta} \psi_{\alpha}).
\]

5. The associated variational problem and harmonic superfunctions

Next, we want to study the Euler-Lagrange equations corresponding to the Berezinian problem defined by \( L \). To this end, we recall that \( G \) determines a global section \( \xi_G \) of the Berezinian sheaf \( \text{Ber}(\mathcal{A}) \), which is called the Riemannian Berezinian volume element associated to \( G \). In coordinates:

\[
\xi_G = \left[ d^G x^1 \wedge \cdots \wedge d^G x^m \otimes \frac{\partial}{\partial x^{-1}} \circ \cdots \circ \frac{\partial}{\partial x^{-n}} \right] |G|
\]
where $|G| = \sqrt{\text{Ber}(G_{\alpha\beta})}$ and $\text{Ber}$ is the Berezinian determinant (note that $|G|$ is even and it only depends on the even coordinates $x^i$). The Euler-Lagrange equations are

\[
\begin{align*}
\frac{\partial \lambda}{\partial t} - \frac{d}{dx^i} \frac{\partial \lambda}{\partial t_i} & - \frac{d}{dx^{-j}} \frac{\partial \lambda}{\partial t^{-j}} = 0, \quad 1 \leq i \leq m, \quad 1 \leq j \leq n \\
\frac{\partial \lambda}{\partial \tau} - \frac{d}{dx^i} \frac{\partial \lambda}{\partial \tau_i} & + \frac{d}{dx^{-j}} \frac{\partial \lambda}{\partial \tau^{-j}} = 0, \quad 1 \leq i \leq m, \quad 1 \leq j \leq n,
\end{align*}
\]

with $\lambda = |G| \cdot L$.

Taking the explicit expression (2) or (3) into account, it is readily seen that these equations reduce respectively to

\[
\frac{1}{|G|} \frac{\partial |G|}{\partial x^k} g^{kj} \tau_{-j} + \frac{\partial g^{kj}}{\partial x^{-j}} \tau_{-j} + g^{kj} \tau_{-j} + g^{kj} = 0
\]

\[
\frac{1}{|G|} \frac{\partial |G|}{\partial x^k} g^{kj} \tau_{-j} + \frac{\partial g^{kj}}{\partial x^{-j}} \tau_{-j} + g^{kj} \tau_{-j} + g^{kj} = 0.
\]

Classically, the Euler-Lagrange variational equations for the sigma model with target manifold $\mathbb{R}$ are those of the harmonic functions, characterized by $\Delta f = 0$, where $\Delta$ is the ordinary Laplacian: $\Delta f = \text{div}(\text{grad} f)$ (see [16,15]). We would like to show that this is still true in the graded setting, that is, that equations (1) when evaluated on sections are the local expression of the harmonic superfunctions, aside from constant factors. Of course, to define the super-Laplacian $\Delta$ one must specify first what is meant by div and grad in a supermanifold.

If $f \in \wedge \mathcal{E}$ is a superfunction on $(M, \wedge \mathcal{E})$ with graded metric $G$, its gradient is defined as the supervector field given by $\langle \text{grad} f, D \rangle$, for all $D \in \text{Der}(\wedge \mathcal{E})$.

**Proposition 4.** For a superfunction $f \in \wedge \mathcal{E}$, the following local expression holds true,

\[
\text{grad} f = g^{ij} \frac{\partial f}{\partial x^{-j}} \frac{\partial}{\partial x^i} + g^{kl} \frac{\partial f}{\partial x^k} \frac{\partial}{\partial x^{-l}}.
\]

**Proof.** Indeed, making use of the $\wedge \mathcal{E}$-bilinearity of $G$ and the explicit form (3), if $\text{grad} f = A^\alpha \frac{\partial}{\partial x^\alpha}$, then

\[
\frac{\partial f}{\partial x^{-j}} = \left< \text{grad} f, \frac{\partial}{\partial x^{-j}} \right> = A^\alpha \left< \frac{\partial}{\partial x^\alpha}, \frac{\partial}{\partial x^{-j}} \right> = A^i g_{ij}.
\]

Hence $A^i = g^{ij} \frac{\partial f}{\partial x^{-j}}$, and similarly for $A^{-l} = g^{kl} \frac{\partial f}{\partial x^l}$. \qed

Next, we study the divergence. If we consider the Riemannian Berezinian $\xi = \xi|G|$ and compute the divergence of $D \in \text{Der}(\wedge \mathcal{E})$ with respect to it, we obtain the following (recall the properties of the Lie derivative of sections of the Berezinian sheaf from section 3):

**Proposition 5.** For any $D \in \text{Der}(\wedge \mathcal{E})$, we have

\[
\text{div} D = (-1)^{|D||\xi|} \frac{1}{|G|} \frac{\partial}{\partial x^\alpha} (|G| \cdot D^\alpha).
\]
Proof. It goes as follows:

\[ L^G_D \xi_G = L^G_D \cdot \xi \cdot |G| \]

\[ = (-1)^{v_\xi} L^G_D \cdot (|G| \cdot D^\alpha) \]

\[ = (-1)^{v_\xi} \left( L^G_{\partial/\partial x^\alpha} (|G| \cdot |D^\alpha|) + (-1)^{v_\xi} \xi \cdot \frac{\partial}{\partial x^\alpha} (|G| \cdot D^\alpha) \right) \]

\[ = \xi \cdot \frac{\partial}{\partial x^\alpha} (|G| \cdot D^\alpha) \]

\[ = \xi \cdot |G| \cdot \frac{1}{|G|} \frac{\partial}{\partial x^\alpha} (|G| \cdot D^\alpha) \]

\[ = \xi_G \cdot \frac{1}{|G|} \frac{\partial}{\partial x^\alpha} (|G| \cdot D^\alpha). \]

Finally, from (2) and (3), we obtain the expression for the super-Laplacian of a superfunction \( f \in \wedge E \), namely,

**Proposition 6.** For any superfunction \( f \in \wedge E \), the equation \( \Delta(f) = 0 \) is equivalent to:

\[ \frac{1}{|G|} \frac{\partial}{\partial x^i} g^{ij} f_{-j} + \frac{\partial g^{ij}}{\partial x^i} f_{-j} + g^{ij} f_{i,-j} + g^{ij} f_{-i,j} = 0. \]

Proof. By a direct computation:

\[ (-1)^{(|f|+1)\xi} \Delta(f) = (-1)^{(|f|+1)\xi} \text{div} \text{grad} f \]

\[ = \frac{1}{|G|} \frac{\partial}{\partial x^i} (|G| \cdot \text{grad} f) \]

\[ = \frac{1}{|G|} \left( \frac{\partial}{\partial x^i} \left( |G| \cdot g^{ij} \frac{\partial f}{\partial x^{-j}} \right) + \frac{\partial}{\partial x^i} \left( |G| \cdot g^{ik} \frac{\partial f}{\partial x^{-k}} \right) \right) \]

\[ = \frac{1}{|G|} \frac{\partial}{\partial x^i} g^{ij} f_{-j} + \frac{\partial g^{ij}}{\partial x^i} f_{-j} + g^{ij} f_{i,-j} + g^{ij} f_{-i,j}. \]

Separating the even and odd factors, we obtain (1) precisely when evaluated on sections of \( p: \mathbb{R}^{1|1} \times (M, \wedge E) \rightarrow (M, \wedge E) \). We have thus achieved the following characterization:

**Theorem 7.** The harmonic superfunctions are precisely the solutions to the Euler-Lagrange equations of the \((1|1)\)-supersymmetric sigma model.

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