THE EXTREME OVERABUNDANCE OF MOLYBDENUM IN TWO METAL-POOR STARS

RUTH C. PETERSON
Astrophysical Advances and UCO/Lick, 601 Marion Pl, Palo Alto, CA 94301
Received 2011 May 23; accepted 2011 August 10; published 2011 November 2

ABSTRACT

We report determinations of the molybdenum abundances in five mildly to extremely metal-poor turnoff stars using five Mo II lines near 2000 Å. In two of the stars, the abundance of molybdenum is found to be extremely enhanced, as high or higher than the neighboring even-Z elements ruthenium and zirconium. Of the several nucleosynthesis scenarios envisioned for the production of nuclei in this mass range in the oldest stars, a high-entropy wind operating in a core-collapse supernova seems uniquely capable of the twin aspects of a high molybdenum overproduction confined to a narrow mass range. Whatever the details of the nucleosynthesis mechanism, however, this unusual excess suggests that very few individual nucleosynthesis events were responsible for the synthesis of the light trans-Fe heavy elements in these cases, an unexpected result given that both are only moderately metal-poor.

Key words: nuclear reactions, nucleosynthesis, abundances – stars: abundances – stars: individual (HD 76932, HD 94028, HD 140283, HD 160617, HD 211998) – stars: Population II – ultraviolet: stars

1. INTRODUCTION

The relative abundances of elements heavier than iron in low-mass stars of low iron abundance bear silent witness to the exploding massive stars in which their metal content was created. Because old massive stars have long since evolved, their solar-mass counterparts are the only surviving stellar relics in which these events have been recorded. Their heavy-element abundance distributions, reviewed by Sneden et al. (2008), can yield critical diagnostics of the objects and environments that formed the material, and the sequences of events that resulted in the buildup of the halo and disk of our Galaxy.

In single stars of metallicity below one-thirtieth solar, [Fe/H] < −1.5, the heavy elements from barium (Z = 56) onward owe their existence to the r-process (rapid neutron capture on seed iron nuclei). Their proportions with respect to one another are the same in all metal-poor stars, despite the wide range of two orders of magnitude observed in their overall abundance with respect to iron. The most probable r-process site is Type II supernovae, whose progenitors are short-lived massive stars. In more metal-rich single stars, [Fe/H] ≥ −1.5, elements begin to appear that are created by the s-process (slow neutron capture), in pulsations in intermediate-mass asymptotic giant branch (AGB) stars. The AGB evolutionary time of a few 100 Myr suggests a time delay of this order in star formation at and above this metallicity.

More complex is the origin of the lightest trans-Fe elements gallium through cadmium (Z = 31 to 48). These elements have been attributed in varying degrees to the p-process (proton capture; Burbidge et al. 1957; Arnould 1976; Woosley & Howard 1978), a “weak” s-process (Clayton 1968; Kappeler et al. 1989), a “weak” r-process (Seeger et al. 1965; Cowan et al. 1991), a specific “light element primary process” (LEPP; Travaglio et al. 2004) such as the vp process of Fröhlich et al. (2006), and/or the low-entropy domain of a neutrino wind above the neutron star formed in a Type II supernova (e.g., Freiburghaus et al. 1999).

In this paper, we establish and discuss the abundances of the two light trans-Fe elements molybdenum and ruthenium (Mo, Ru; Z = 42, 44) in five metal-poor stars whose enhancements of heavy r-process elements are mild. In two of the five stars, HD 94028 and HD 160617, the molybdenum abundance is extremely elevated, with important ramifications for the synthesis process and the number of synthesis events contributing to the light trans-Fe elemental abundances.

2. CURRENT OBSERVATIONAL CONSTRAINTS ON LIGHT TRANS-Fe ELEMENTS

As reviewed by Lodders (2010), stringent constraints of nucleosynthesis models are derived from isotopic abundances of meteorites (e.g., Pellin et al. 2006). These reflect the integral of the products of all processes incorporated into the pre-solar nebula—a single detailed example in space and time.

Isotopes of the light trans-Fe elements have proven to be among the most difficult to reproduce. Molybdenum is especially problematical; its solar system p-process isotopic fraction of ∼25% is larger than that of any other trans-Fe element (Lodders 2010). This is a stumbling block for models invoking the s-process in low-mass AGB stars (e.g., Lugano et al. 2003). Such a scenario is unlikely to apply in any case to the low-metallicity stars of the halo, as it relies on the preexistence of AGB stars of solar metallicity and also of quite low mass, 1.5 M⊙, with accordingly long main-sequence lifetimes. In contrast, Hoffman et al. (1996) succeeded in directly producing light p-process nuclei with specific choices of entropy S and electron fraction Ye in a neutrino-driven wind. They noted that this is a primary process, one in which “the r-process and some light p-process nuclei may be coproduced.”

Recently, Faruqui et al. (2009) reproduced all seven of the solar isotopes of molybdenum by selecting models from a parameterized grid of calculations based on a high-entropy wind (HEW) operating in Type II supernovae. They find it “can coproduce the light p-, s-, and r-process isotopes between Zn (Z = 30) and Ru (Z = 44) at electron abundances in the range 0.450 ≤ Ye ≤ 0.498 and low entropies of S ≤ 100–150. Under these conditions, the light trans-Fe elements are produced in a charged-particle (αp) process, including all p-nuclei up to 96,98Ru. In our model, no initial SS (solar system), s- or r-process seed composition is invoked; hence, this nucleosynthesis component is primary.” In part because “the overall yields of the light trans-Fe elements decrease with increasing Ye,” they conclude that “more quantitative answers to questions concerning the astrophysical site of the compositions of the...
LEPP elements between Sr (Z = 38) and Cd (Z = 48), as well as all of the n-capture elements, will require more and higher quality observational data and also more realistic values of entropy superpositions derived from hydrodynamical models. 

Abundances of lighter and heavier elements in metal-poor stars are already providing further constraints. Roederer et al. (2010a) derived abundances for zinc, yttrium, lanthanum, europium, and lead (Zn, Y, La, Eu, and Pb; Z ≥ 30, 39, 57, 63, and 82) in 161 metal-poor stars with [Fe/H] < −1.4. Based on models of the s-process in AGB stars, they used [Pa/Fe] to identify stars with no discernible s-process contribution, and concluded that s-process elements were largely absent from progenitor material at these low metallicities. Because a scatter remained in [La/Fe] in those stars with relatively low r-process content, they confirmed the result emphasized earlier by Honda et al. (2007) that the ratio of light r-process to heavy s-process elements varied widely among metal-poor stars. Roederer et al. (2010a) also confirmed an anti-correlation between Y/Eu and Eu/Fe (François et al. 2007), and showed that Y production was decoupled from both Zn and Fe. They were able to reproduce the range of Y/Eu ratios with simulations of HEW models that explore the effects of a range of entropies (Farouqi et al. 2010).

Molybdenum itself, previous abundance determinations in metal-poor stars are restricted to giants and subgiants, in which near-UV and optical Mo I lines are detectable. Except for giants with extreme r-process enhancements (e.g., Sneden et al. 2003), published [Mo/Fe] values are all near solar. Table 1 summarizes values for [Mo/Fe], and other light trans-Fe elements where available, for eight field halo stars with [Fe/H] ≤ −1.4 and [Eu/Fe] < 0.9, as an indicator of mild r-process enhancement. Of the 16 globular cluster studies listed in Table 1 of Roederer (2011), only one presents results for Mo. In that study, for eight giants in the globular cluster M5, Lai et al. (2011) find [Fe/H] = −1.43, [Eu/Fe] = +0.49, [Zr/Fe] = +0.34, and [Mo/Fe] = −0.10, with no significant star-to-star variation. The referee adds that Yong et al. (2008) have derived [Mo/Fe] for 11 stars in M4 and two in M5, none of which has [Mo/Fe] > +0.4. Since M4 has a subgroup of stars with significant s-process contribution (Figure 1, panel 4 of Roederer 2011), its non-s-process [Mo/Fe] upper bound may even be lower. Among the dozen normal field and cluster giants in which molybdenum has been studied to date, then, none has [Mo/Fe] > +0.4.

| Star   | Wavelength (Å) | Instrument | Program     | Date (UT) | Time (ks) | Reduction |
|--------|----------------|------------|-------------|------------|-----------|-----------|
| HD 140283 | 1950–2300 | STIS E230H | GO 7348 | 1999 Apr 9 | 18.32 | StarCat uvsum2126 |
| 2378–2891 | STIS E230H | GO 9455 | 2002 Aug 22 | 5.28 | IRAF |
| 2885–3147 | STIS E230H | GO 9491 | 2003 Jul 11,12,13,16,17,22,23,24 | 62.57 | StarCat 52823–52844 |
| 3080–5953 | HIRES | U35H | 2005 Mar 17 | 0.60 | HiRedux |
| HD 160617 | 1880–2150 | STIS 230H | GO 8197 | 1999 Oct 29,30; 2000 Mar 15,16; Aug 31 | 39.39 | StarCat 51480–51787 |
| 3057–3873 | UVES | 65L-0507(A) | 2000 Apr 9 | 3.00 | Pipeline |
| 4400–6780 | HIRES | Hs6H | 2000 May 28 | 0.42 | Extracted |
| HD 94028 | 1880–2150 | STIS 230H | GO 8197 | 2000 May 16,21,26,21 | 33.05 | IRAF |
| 2278–3120 | STIS 230M | GO 7402 | 1998 Dec 18 | 0.60 | IRAF |
| 3050–4989 | UVES | 072.B-0585(A) | 2004 Mar 10 | 0.75 | NGSL |
| HD 76932 | 1880–2150 | STIS 230H | GO 9804 | 2004 Feb 19,21 | 23.86 | StarCat 53054–53056 |
| 3022–4975 | UVES | 266.D-5655(A) | 2001 Mar 14 | 0.34 | Pipeline |
| HD 211998 | 1880–2150 | STIS 230H | GO 9804 | 2004 Aug 26,27 | 29.40 | IRAF |
| 3040–1040 | UVES | 266.D-5655(A) | 2002 Feb 9 | 0.60 | Pipeline |

3. STELLAR SPECTRA

In this work we provide additional support and constraints for HEW production of light trans-Fe elements, by determining the abundances of Mo and Ru in metal-poor turnoff stars from Mo II and Ru II lines near 2000 Å in high-resolution spectra taken with the Space Telescope Imaging Spectrograph (STIS). Five such spectra were found in the MAST archive, the Multimission Archive at the Space Telescope Science Institute (STScI). To constrain the molybdenum abundance scale and to derive abundances for other elements, archival near-UV and optical echelle spectra were analyzed for the same stars.

Table 2 lists for each spectral region the spectra employed for each star. Reductions by others were adopted from StarCat.
(Ayres 2010), the UVES pipeline,1 Keck HIRES archival extractions, and the UVES ground-based spectral programs of the Next Generation Spectral Library (Gregg et al. 2004). Our own reductions used the IRAF2 environment. We performed bias and dark removal, co-added multiple spectral images of the same object with cosmic-ray removal, extracted orders with removal of sky and local interorder background, corrected the dispersion using Th–Ar exposures, and rectified the continuum and spliced together adjacent orders.

4. SYNTHETIC SPECTRAL ANALYSIS

We have derived stellar parameters and abundances by matching each stellar spectral observation to theoretical spectra calculated for each star using an updated version of the SYNTHE program of Kurucz (1993b). We input a list of molecular and atomic line transitions with wavelengths, energy levels, and gf-values, and a model atmosphere characterized by effective temperature $T_{\text{eff}}$, surface gravity $\log g$, microturbulent velocity $V_t$, and logarithmic iron-to-hydrogen ratio $[^{\text{Fe}}/H]$ with respect to that of the Sun. Our models are interpolated in the grid of Castelli & Kurucz (2003).

The line lists are based on the Kurucz gfhy3 atomic lines with known energy levels (“laboratory” lines), along with Kurucz diatomic molecular line lists4 and TiO lines5 (Schwenke 1998). We have modified these extensively in the near-UV and optical, 2200–9000 Å, by matching calculations to echelle spectra of standard stars. Starting with weak-lined stars, we calculated each spectrum, adjusted gf-values singly for atomic lines and as a function of band and energy for molecular lines, and guessed identifications of “missing” lines, those present in the spectra but not in the laboratory line list, which become an extreme problem in the UV. Peterson et al. (2001) detail the procedure, and that work and Peterson (2008) show that our calculations agree well with observed optical and mid-UV spectra of nearby mildly metal-poor stars.

Following these procedures, we continued the line modifications into the 2000 Å region. This better defines the local continuum, especially in the two stronger-lined stars HD 76932 and HD 211998. All missing lines were simply assumed to be Fe i lines with lower excitation potential below 1 eV, as these are extremely common among the identified lines. We added these at wavelengths where the spectra of HD 76932 and HD 211998 showed absorption at the same wavelength that was not matched by any laboratory line in the Kurucz atomic line lists, even when its $gf$-value is increased by as much as 2.5 dex.

We find stellar parameters strictly from the spectra, and not from colors. Effective temperature $T_{\text{eff}}$ comes from demanding that the same abundance emerge from low- and high-excitation lines of same species (usually Fe i). Gravity log $g$ comes from the wings of other strong lines. Demanding no trend in abundance with line strength sets microturbulent velocity $V_t$. Iron abundance $[^{\text{Fe}}/H]$ follows by matching relatively unblended weak lines, as do the abundances of other elements. The resulting uncertainties are typically 0.1–0.2 dex in $[^{\text{X}}/\text{Fe}]$ for element X, if represented by at least three lines whose blending, if any, is reliably modeled, and whose $gf$-values are well determined. Comparing Fe i and Fe ii abundances confirms or refines gravity, and the wings of the Balmer lines confirm $T_{\text{eff}}$. These agree with other $T_{\text{eff}}$ diagnostics only when convection overshoot is turned off, as is true of the Castelli & Kurucz (2003) models but not those of Kurucz (1993a). Table 3 lists the resulting stellar model parameters and abundances.

Table 3

| Star         | $T_{\text{eff}}$ | $\log g$ | $[^{\text{Fe}}/H]$ | $[^{\text{Eu}}/\text{Fe}]$ | $[^{\text{Y}}/\text{Fe}]$ | $[^{\text{Zr}}/\text{Fe}]$ | $[^{\text{Mo}}/\text{Fe}]$ | $[^{\text{Ru}}/\text{Fe}]$ | $[^{\text{La}}/\text{Fe}]$ |
|--------------|-----------------|----------|---------------------|-----------------------------|----------------------------|----------------------------|----------------------------|----------------------------|-----------------------------|
| HD 140283    | 5700            | 3.6      | −2.6                | < −0.9                      | −0.4                       | −0.1                       | +0.2                       | < +1.0                     | ...                         |
| HD 160617    | 6000            | 3.8      | −1.8                | 1.2                         | +0.6                       | +0.0                       | +0.4                       | +0.8                       | +0.6                       |
| HD 94028     | 6050            | 4.5      | −1.4                | −1.2                        | +0.3                       | +0.2                       | +0.5                       | +1.0                       | +0.7                       |
| HD 76932     | 5900            | 4.1      | −0.9                | 1.2                         | +0.4                       | +0.0                       | +0.2                       | +0.6                       | +0.4                       |
| HD 211998    | 5300            | 3.1      | −1.5                | 1.5                         | +0.2                       | +0.2                       | +0.5                       | +0.5                       | −0.06                       |

Note. Units: $T_{\text{eff}}$, K, $\log g$, km s$^{-1}$.

We find stellar parameters strictly from the spectra, and not from colors. Effective temperature $T_{\text{eff}}$ comes from demanding that the same abundance emerge from low- and high-excitation lines of same species (usually Fe i). Gravity log $g$ comes from the wings of other strong lines. Demanding no trend in abundance with line strength sets microturbulent velocity $V_t$. Iron abundance $[^{\text{Fe}}/H]$ follows by matching relatively unblended weak lines, as do the abundances of other elements. The resulting uncertainties are typically 0.1–0.2 dex in $[^{\text{X}}/\text{Fe}]$ for element X, if represented by at least three lines whose blending, if any, is reliably modeled, and whose $gf$-values are well determined. Comparing Fe i and Fe ii abundances confirms or refines gravity, and the wings of the Balmer lines confirm $T_{\text{eff}}$. These agree with other $T_{\text{eff}}$ diagnostics only when convection overshoot is turned off, as is true of the Castelli & Kurucz (2003) models but not those of Kurucz (1993a). Table 3 lists the resulting stellar model parameters and abundances.

5. MOLYBDENUM AND RUTHENIUM ABUNDANCES

As seen in these figures, the calculated spectra match the fitted spectra quite well. However, the abundance determinations are affected by systematic errors, due to the uncertainties in continuum placement and in $gf$-values. These are difficult to judge.

Continuum placement is affected by line absorption, whose modeling depends on $gf$-value adjustment and on correct assignment of the wavelength, species, and lower excitation potential of missing lines. The number of missing lines grows dramatically toward the UV, and consequently the continuum becomes less well defined below 2000 Å, even for the weak-lined star HD 160617. This is illustrated in the top left panel in Figure 1. Line-strength adjustment was attempted only over 1887.6–1889.0 Å, as spectral data are currently lacking for the weakest-lined star HD 140283, and many significant lines are missing. The latter shortcoming is likely to be alleviated by revising the input line list to include recent recalculations by Kurucz (2011), posted on the Kurucz Web site in 2011 April, which include 10 times as many lines of Mg, Si, Ca, and the iron-peak elements.

Their inclusion might change the Ru abundances of HD 160617 and HD 94028 by up to 0.3 dex. While the blending of the Ru i line itself can be modeled primarily with identi-

1 http://www.eso.org/sci/software/pipelines/
2 IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.
3 http://kurucz.harvard.edu/LINELISTS/GFHYPER100/
4 http://kurucz.harvard.edu/LINELISTS/LINESMOL/
5 http://kurucz.harvard.edu/molecules/TiO/
Figure 1. Comparisons are shown between observed and calculated spectra in nine spectral regions, indicated by wavelength in Ångstroms at the bottom. Plots for the five individual stars are offset vertically; ticks on the y-axis indicate one-tenth of the normalized continuum level. The HD number of each star is given above its plot. The heavy line is its observed spectrum, and the light line its calculated spectrum. The strongest lines in the calculated spectrum are identified at the top. First are the digits following the decimal place of the line center wavelength (in vacuum for the bluest region and in air otherwise). Next is given the species giving rise to the line; a colon indicates a “missing” line whose identification was assumed to be Fe i. Following this are the lower excitation potential of the line in eV, an indicator of its strength, and its log $g_f$-value. Three calculations are shown near Mo and Ru lines; these are expanded and described in Figure 2.
Figure 2. Observed and calculated spectra of Figure 1 are shown in the same nine spectral regions, and with the same calculations and labels, but on an expanded wavelength scale in the vicinity of the Mo and Ru lines. The middle light line is calculated assuming the abundances of Table 3. The exception is the choice of [Ru/Fe] = +0.1 ±0.3 for HD 140283, as no Ru i nor Ru ii lines with reliable gf-values were detected for this star. The weaker and stronger light lines in each plot indicate calculations with Mo and Ru abundances 0.3 dex lower and higher. Red arcs highlight a few cases in which the line indicated was the only line calculated.

fied lines, if missing lines are also present but currently unrecognized, the currently inferred ruthenium abundance may be overestimated. However, if there are more missing lines in the adjacent regions used to set the continuum, its level is underestimated, and the inferred ruthenium abundance may be underestimated.

Scale errors can occur in both laboratory and theoretical gf-values. We examined these wherever possible. To calculate the
UV Mo II lines, we increased by 0.133 dex the Mo II log $gf$-values of Sikström et al. (2001) for the five Mo II transitions that originate from the ground state. This is the difference between their value and the recent result of Lundberg et al. (2010) for the 2082 Å Mo II line, the only ground-state Mo II line in common between the two studies. We followed Evans et al. (2006) in adopting the Mo I $gf$-values of Whaling & Braeul (1988) and the Ru I results of Wickliffe et al. (1994). For Ru II, Johansson et al. (1994) provide 18 experimental $gf$-values, and Palmeri et al. (2009) give theoretical values. Both wavelengths and $gf$-values proved unreliable for ruthenium lines with data solely from other sources.

For HD 76932 and HD 211998, the molybdenum abundances were determined from the optical Mo II line at 3864 Å. For HD 94028, we adopted the molybdenum abundance which was just consistent with the non-detection of the Mo I 3797 Å line. For HD 160617 and HD 140283, the molybdenum abundances were determined by matching Mo II line strengths, as these lines are rather weak, and Mo I is not detected. The Mo I and Mo II $gf$-value scales thus appear to be on a consistent scale. This is encouraging, as the Mo I $gf$-values have yielded a solar molybdenum abundance (Biernot et al. 1983) within 10% of the meteoritic value (Lodders 2010).

For ruthenium, we sought the strongest Ru I lines in the optical and Ru II lines in the UV. It is again encouraging that Ru I and Ru II lines give consistent results for HD 76932 and HD 211998. For the other stars, however, the ruthenium lines with $gf$-values from the three sources above that lie within the wavelength ranges observed at high resolution are weak and blended. In HD 160617 and HD 94028, our Ru values are based on these plots of the two Ru II lines and the Ru I line at 3499.942 Å. For HD 140283, the limit is set by calculations not shown, adopting [Ru/Fe] $= +1.0$ for the Ru II 2102.307 Å line.

6. CADMIUM

Cadmium ($Z = 48$) is currently constrained in only two stars. In HD 94028, [Cd/Fe] $\sim 0.0$ is found from lower-resolution E230M spectra of the Cd I 2288.018 Å line, adopting log $gf = +0.15$ dex (Roeder et al. 2010b). The HD 140283 E230H spectra show artifacts in this region; data are lacking entirely for other stars. In HD 140283, [Cd/Fe] $< 0.0$ is inferred from the Cd II 2144.393 Å line, even with log $gf = -0.11$ (Andersen & Soerenzen 1973), which is lower than more recent $gf$-values for this line ($-0.04$ to $-0.12$: Xu et al. 2004; Mayo et al. 2005). In the other stars, this line is possibly blended. Obtaining E230H data for the Cd I line at 2288.018 Å for these stars should pinpoint cadmium abundances, by resolving its own potential blends, determining [Cd/Fe] in HD 140283, and so fixing the Cd II 2144.393 Å $gf$-value, and then using Cd I and Cd II together to set [Cd/Fe]. This would more strongly constrain the extent of light trans-Fe elemental overabundances in atomic number $Z$ and therefore the models of their production.

7. THE S-PROCESS CONTRIBUTION TO STELLAR HEAVY-ELEMENT ABUNDANCES

The $s$-process may contribute to the Ru and Mo abundances, especially for the more metal-rich stars. To check this we have included [La/Fe] abundances from the 3995.75 Å and 4086.71 Å La II lines, with $gf$-values from Lawler et al. (2001) and hyperfine splitting structure from Evans et al. (2006). We compare our Table 3 values with the trend of [La/Fe] versus [Eu/Fe] for $r$-only field halo stars in the fourth panel of Figure 1 of Roederer (2011). HD 76932 and HD 211998 lie within 0.05 dex of the mean relation. HD 160617 and HD 94028 fall $< 0.3$ dex above the line, marginally beyond the extent of the $r$-only field stars in that figure. We conclude that any excess $s$-process contribution is small in these four stars.

For HD 140283, none of these La II lines were detected, but at [Fe/H] $= -2.6$ an $s$-process contribution is unlikely (Roeder et al. 2010a). Gallagher et al. (2010) find ambiguous results for its $s$-process content from fitting the 4554 Å and 4934 Å Ba II line profiles. That work and ours agree that [Eu/Fe] $< -0.9$ in this star. Its abundance pattern strongly resembles that of HD 88609 and HD 122563 in showing larger deficiencies of heavy than light trans-Fe elements, with low relative abundances overall (Honda et al. 2007).

8. RESULTS AND IMPLICATIONS FOR NUCLEOSYNTHESIS

Thus we conclude that in HD 94028, molybdenum is extremely enhanced, more so than ruthenium. Mo is also highly enhanced in HD 160617. Both of these stars, indeed four of the five metal-poor turnoff stars studied here, have higher [Mo/Fe] values than any of the eight metal-poor giants with mild (if any) $r$-process enhancements listed in Table 1. Yet none has [La/Fe] enhancements more than 0.3 dex higher than expected from its [Eu/Fe] value, based on Figure 1 of Roederer (2011). The total range in [Mo/Fe] among these five stars is 0.8 dex, while the total range in [La/Fe] is 0.4 dex. Among these stars with modest or no $r$-process enhancements, no correlation is seen between [Eu/Fe] and [Mo/Fe]. The extreme enhancements are rather narrowly confined to Mo, diminishing toward Zr and Y and toward Ru and beyond.

Since both $s$- and $r$-process nucleosynthesis tend to produce similar enhancements over a range of non-magic neighboring even-Z elements, the production of molybdenum and ruthenium in HD 94028 most probably involves another process. As noted above, a HEW is capable of overproducing light trans-Fe elements if the wind parameters are right. Faruqui et al. (2010) express this in terms of $Y_e$: $Y_e = 0.498$ yields Sr, Y, Zr, and Nb ($Z = 38-41$); $Y_e = 0.496$ yields Mo and Ru; $Y_e = 0.490$ yields Rh, Pd, and Ag ($Z = 45-47$); and $Y_e = 0.482$ yields Cd and beyond ($Z \geq 48$). Clearly, a HEW with a limited parameter range seems able to reproduce the strong excess of molybdenum with less strong excesses at zirconium and beyond ruthenium.

These unique factors also suggest that very few individual nucleosynthesis events were incorporated into the stars with extreme molybdenum abundances. This is especially remarkable for HD 94028, given its rather high metallicity, [Fe/H] $= -1.4$. From Table 7 of Faruqui et al. (2010), the yield from an individual HEW event with $Y_e = 0.496$ is $10^{-5} M_\odot$; this is more than adequate. However, the narrow entropy range means that multiple HEW events with a range of parameters must be avoided during the buildup of the iron abundance.

The referee has pointed to another example, from Aoki et al. (2007). They find that the star COS 82 in the dwarf galaxy Ursa Minor, with [Fe/H] $= -1.5$, has a very high heavy element enhancement with an $r$-process signature. They note that such high $r$-process enhancements are found in Galactic stars only below [Fe/H] $= -2.5$ and suggest that “the neutron-capture elements of COS 82 might be provided by a single event.”

Such a scenario has recently become theoretically more feasible. Bland-Hawthorn et al. (2011) have found that a dark matter halo of about $3 \times 10^6 M_\odot$ in gas, a lower-mass
system than previously thought, can still form stars from retained products released by an exploding supernova. The lower mass results from the reduced sweeping of products when clumpiness of the medium and off-center supernovae explosions are included.

However, as discussed above, we know of no prior evidence for single-supernova production of any group of trans-Fe elements in moderately metal-poor Galactic halo stars. More Mo abundances in halo stars with \([\text{Fe}/\text{H}] \leq -1.4\) are needed, to verify rarity and define the frequency of occurrence of high \([\text{Mo}/\text{Fe}]\) values as a function of metallicity. A survey that includes stars with well-established kinematics might reveal whether high-Mo stars might have been formed in captured/dissipated dwarf galaxies. For stronger-lined turnoff stars like HD 76932 and HD 211998, either cool or of moderate metallicity, Mo and Ru can be determined from archival ground-based spectra, as shown in Figure 1. For the most metal-poor r-normal stars, however, the UV lines are required.

Determination of the abundances of more light trans-Fe elements in HD 94028 and HD 160617 seems critical as well to provide stronger HEW support and constraints for parameterized modeling. The steeper the falloff of abundance enhancements from Mo to Ru and beyond, the more narrow the entropy range implied in HEW synthesis. Obtaining high-resolution spectra redward of 2200 Å for these stars might reveal Nb and Pd abundances, as well as pin down Ru and Cd.

Both of these efforts are vital in revealing whether the very high molybdenum enhancements reported here are unique. This in turn would help establish to what extent they truly are the result of very few events in some stars, even at quite high metallicity.

We thank M. Spitze for insightful discussions, J. Bland-Hawthorn for describing his recent work, J. X. Prochaska for providing his reductions of the Keck HIRES data, and D. Silva and R. Hanuschik for providing the reduced UVES NGSL spectra. We also appreciate the helpful report of the referee. Ground-based spectra are based on observations made with ESO Telescopes at the Paranal Observatory with the UVES spectrograph, as shown in Figure 1. For the most metal-poor r-normal stars, however, the UV lines are required.

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