Arctic curves in path models from the tangent method

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Abstract
Recently, Colomo and Sportiello introduced a powerful method, known as the tangent method, for computing the arctic curve in statistical models which have a (non- or weakly-) intersecting lattice path formulation. We apply the tangent method to compute arctic curves in various models: the domino tiling of the Aztec diamond for which we recover the celebrated arctic circle; a model of Dyck paths equivalent to the rhombus tiling of a half-hexagon for which we find an arctic half-ellipse; another rhombus tiling model with an arctic parabola; the vertically symmetric alternating sign matrices, where we find the same arctic curve as for unconstrained alternating sign matrices. The latter case involves lattice paths that are non-intersecting but that are allowed to have osculating contact points, for which the tangent method was argued to still apply. For each problem we estimate the large size asymptotics of a certain one-point function using LU decomposition of the corresponding Gessel–Viennot matrices, and a reformulation of the result amenable to asymptotic analysis.

Keywords: non-intersecting lattice paths, continuum limit, arctic curve

(Some figures may appear in colour only in the online journal)

1. Introduction

It is now well known that under certain conditions tiling problems of finite plane domains display an ‘Arctic curve’ phenomenon [1, 9] when the size of the domains becomes very large, namely a sharp separation between ‘crystalline’ (i.e. regularly tiled) phases, typically induced by corners of the domain, and ‘liquid’ (i.e. disordered) phases, away from the boundary of the domain. A particular subclass of such problems are the so-called ‘dimer’ models, where tiling configurations are replaced by a dual notion of dimers, i.e. occupation of the edges of a
given domain of a lattice by objects (dimers) in such a way that each vertex of the domain is covered by a unique such object. These have received considerable attention over the years, culminating in general asymptotic results and a characterization of the arctic curves as solving some optimization problem [11–13].

A common denominator between all these models is the existence of a reformulation in terms of non-intersecting lattice paths (NILP). For tilings, it is due to the existence of conservation laws (local properties that propagate throughout the domain), giving rise to configurations of paths (often called De Bruijn lines [5]) in bijection with the tilings. For dimers, it is related to the description of configurations via so-called zig–zag paths [10].

Starting with the pioneering works [1, 9] on domino tilings of the Aztec diamond, the method used to derive the arctic curve involves the asymptotic analysis of some bulk correlation functions, which can become quite involved. Recently, Colomo and Sportiello [4] came up with a novel approach to the determination of arctic curves, coined the ‘Tangent Method’. Its main advantage is to use the properties of the model at hand in a minimal way, by only requiring computation of boundary correlation functions, a much easier task. The method is based on the non-intersecting lattice path formulation. The idea is to modify the lattice path configurations of a certain domain so as to impose that a single border path escapes and reaches a distant target point outside of the domain. It is argued that for large size this path should leave the arctic curve at some point. Away from the other paths, and for large size, the latter is most likely to follow a straight line between the point where it leaves the arctic curve and the distant target. This line is moreover argued to be tangent to the actual arctic curve. The tangent method consists of determining for a fixed target, the most likely boundary point of the domain at which the escaping path leaves the domain, thus producing a parametric family of lines that are tangent to the arctic curve. The latter is recovered as the envelope of the family of lines obtained by moving the target, say along a line. The intriguing feature of the method is that it seems to extend beyond ordinary non-intersecting lattice path models. In particular, in [4] the authors argue that the method also applies to the case of so-called osculating paths such as those in bijection with configurations of the 6 vertex model or alternating sign matrices (ASMs) at the ice point [16, 19]. As a consequence they obtain an alternative derivation of the ASMs arctic curve providing the same result as earlier calculations based on assumptions of a very different nature [2, 3].

The tangent method however is not rigorous as it relies on some fundamental (yet reasonable) assumptions. On the other hand, the dependence of the shape of the arctic curve on particular choices of boundary conditions is still not well understood. Finally the range of applicability of the method is not known precisely, and the arctic curve computation for ASM by Colomo and Sportiello [4] seems to indicate that it applies to interacting models as well. To tackle these problems, this paper addresses some concrete examples of path/tiling problems and applies the tangent method to determine arctic curves, in order to either (i) compare them to known results (such as the arctic circle of the domino tilings of the Aztec diamond) or (ii) derive some new ones. This motivates our choice of examples below.

We first treat the case of domino tilings of the Aztec diamond for which we recover the celebrated arctic circle result of [1]. This is used as a pedagogical introduction to the tangent method, by first reformulating the tiling problem as a non-intersecting lattice path problem, and by calculating some boundary one-point function of the model. To derive this function, we develop a general approach based on the LU factorization of the Gessel–Viennot matrix for the path model, which is considerably simpler than the standard use of the Kasteleyn matrix machinery to derive bulk correlation functions.

Then we go on to study two lattice path models, both attached to rhombus tilings of particular domains of the triangular lattice. These are chosen because of the specific type of
boundary conditions on the associated tiling problems. They also turn out to be solvable by
the \textit{same} LU factorization method as for the Aztec diamond tilings, giving rise to new explicit
exact expressions for the boundary one-point functions. In the first model, we consider ‘half-
space’ boundary conditions, in which paths are constrained to lie above a fixed line, and for
which we show that the arctic curve is a half-ellipse, such as that obtained in the presence
of very different free boundary conditions for rhombus tilings of a half-hexagon \cite{6}. In the
second, the boundary follows a specific regular pattern and we find that as a consequence the
arctic curve is a parabola. This latter example is meant as a first step in studying more general
boundary conditions with regular patterns, a subject we shall develop in later publications.

Finally we treat the case of vertically symmetric alternating sign matrices (VSASMs), as
an extension of Colomo and Sportiello’s calculation for ASMs, to further probe the tangent
method’s range of applicability. We find that the asymptotic shape of VSASMs is the same
as that of ASMs. This also comes in support of the general reasoning used in \cite{6} to state that
arctic curves for tilings with imposed symmetries which are compatible with the shape of the
domain remain the same when symmetry constraints are lifted.

In our derivations, we have to deal in general with two main complications:

\begin{itemize}
  \item As the tangent method is based on large size estimates, we need to obtain positive expres-
sions for the various counting functions we wish to estimate, however LU decomposition
  naturally yields alternating sum expressions for these. We therefore have to reformulate
  these alternating sums as positive sums, a rather involved procedure.
  \item Each setting for applying the tangent method only yields a portion of the arctic curve,
  hence we must change the setting (and possibly the NILP formulations) to obtain other
  portions.
\end{itemize}

The paper is organized as follows.

In the preliminary section 2, we expose the general tangent method, and how to apply it
in the context of NILP models. In particular we describe a general approach to the computa-
tion of partition functions of NILP and their version with one escaping path, based on the LU
decomposition of the corresponding Gessel–Viennot matrix. We conclude this section with a
review of how to apply generating function methods to this procedure.

Section 3 is devoted to the re-derivation of the arctic circle for domino tilings of the Aztec
diamond by the tangent method applied to their reformulation in terms of (large Schröder)
NILP.

Sections 4 and 5 explore another NILP problem involving non-intersecting Dyck paths
i.e. directed paths on the square lattice that remain in a half-plane. In section 4, we explore
and solve the path problem and derive the naturally associated portion of arctic curve. To get
the entire curve, we first reformulate the path problem as a rhombus tiling problem of a cut
hexagon in section 5, where we derive the rest of the arctic curve by considering an alternative
NILP description of the same tiling configurations. The complete result for the arctic curve is
a half-ellipse inscribed in the cut hexagon.

In section 6, we address yet another NILP problem involving ordinary directed paths on the
square lattice, but with specific constraints on their starting and ending points. For this case we
find that the arctic curve is a portion of parabola.

Section 7 is devoted to the derivation of the arctic curve for VSASMs. We find essentially
the same result as for ordinary ASMs (which do not have the reflection symmetry constraint).

We gather a few concluding remarks in section 8.
2. Preliminaries

2.1. The tangent method

The method we are going to describe was devised by Colomo and Sportiello [4] and applies to a number of problems. First of all tiling problems of plane domains by means of tiles with a few specific shapes and sizes. In many cases, such problems may be reformulated in terms of non-intersecting lattice paths (NILP), for which the tangent method is well-defined. The latter are of course combinatorial problems in their own right, a few of which we will consider in this paper. A more subtle application of the method seems to indicate that it also applies to interacting lattice paths, typically allowed to ‘kiss’ at a vertex, with a particular interaction weight. This is the case for the ‘osculating path’ formulation of the configurations of the six vertex model [4].

Let us consider configurations of families of non-intersecting (directed) lattice paths, say from a set of initial points \( v_0, v_1, \ldots, v_n \) to a set of endpoints \( \tilde{v}_0, \tilde{v}_1, \ldots, \tilde{v}_n \). By lattice paths, we mean paths drawn on a regular directed graph, whose vertices are the lattice points and whose elementary oriented steps are taken along oriented edges of the graph.

The tangent method allows to study the large size asymptotics of those configurations (in a sense defined below). Note that the set of all possible paths from any of the \( v \)’s to any of the \( \tilde{v} \)’s defines a maximal domain \( D \) with a shape depending on the lattice and on the positions of the \( v, \tilde{v} \)’s. By large size asymptotics, we mean the limit of a large (scaled) domain \( D \).

For large size, we expect the configurations to display the following pattern of behavior. As long as the paths remain too close to each other, they will form very regular ‘crystalline’ patterns, whereas far enough from the edges of \( D \) we expect more disorder, i.e. a ‘liquid’ phase. It turns out that in many NILP problems the separation between these phases is very sharp, giving rise asymptotically to separating curves. This is the ‘Arctic Curve Phenomenon’. This can be translated back in terms of tilings by saying that tilings of a very large planar domain \( D \) tend to develop crystalline phases (in particular induced by corners) along the boundaries of \( D \), while a liquid phase develops away from the boundaries. These are separated by arctic curves as well. There is a large literature on this subject originating in the determination of the ‘Arctic Circle’ for the domino tiling of the so-called Aztec diamond [1, 9], and evolving to general descriptions such as that of [11, 13], and further developments studying the fluctuations around the arctic curve as well.

The tangent method allows to predict the precise shape of the arctic curve and is sketched on a particular example in figure 1. It goes as follows. The arctic curve can be explored by considering the most likely outermost paths of the configuration (say paths from \( v_n \) to \( \tilde{v}_n \)), which in the large size limit will tend to portions of the arctic curve. To capture the exact position of the arctic curve, consider another related NILP problem (see figure 1): move the endpoint \( \tilde{v}_n \) to a reachable point \( \tilde{v} \) at a vertex of the underlying lattice graph (e.g. the point \( C \) in figure 1) far away from the other endpoints, in particular outside of \( D \). The corresponding ‘outer path’ \( v_n \to \tilde{v} \) will therefore veer away from the other paths, and its most likely trajectory will be a straight line, when away from the influence of the other paths. We note that this line must be asymptotically tangent to the arctic curve. By moving around the position of the new endpoint \( \tilde{v} \), we may thus generate some parametric family of tangents to the arctic curve, which can be recovered as their envelope.

How to determine the family of tangents? We may record the point \( \tilde{w} \) (e.g. the point \( B \) on figure 1) from which the escaping path \( v_n \to \tilde{v} \) leaves the domain \( D \) (i.e. the last visit to a boundary point of \( D \), and unique vertex from which a step outside of \( D \) is taken). This point is supposed to be already away from the influence of the other paths. So from this point on, the escaping paths will follow the tangent asymptotically, defined as the line through \( \tilde{w} \) and \( \tilde{v} \) (e.g.
the line $BC$ in Figure 1). Assume we can enumerate the NILP configurations from $v_0, v_1, \ldots, v_n$ to $\tilde{v}_0, \tilde{v}_1, \ldots, \tilde{v}_{n-1}, \tilde{v}$, say with partition or counting function $N_{\tilde{v}}$, as well as those from $v_0, v_1, \ldots, v_n$ to $\tilde{v}_0, \tilde{v}_1, \ldots, \tilde{v}_{n-1}, \tilde{w}$, with partition function $Z_{\tilde{w}}$. Then we can perform an asymptotic large size analysis of $N_{\tilde{v}} = \sum_{\tilde{w} \in \partial D} Z_{\tilde{w}} Y_{\tilde{w}, \tilde{v}}$, where $Y_{\tilde{w}, \tilde{v}}$ is the partition function for paths from $\tilde{w}$ to $\tilde{v}$ exiting from $D$ at $\tilde{w}$. Assuming scaling behaviors $\tilde{v} = na, \tilde{w} = nx, n \to \infty, a, x$ two points independent of $n$, we get the following large $n$ leading behavior: $N_{\tilde{w}} \sim \int \! d^2x \, e^{nS(x)}$, for some action functional $S$. By a saddle-point analysis, we find the most likely scaled exit point $x^*$, as a function of $a$. This gives a parametric family of tangents, i.e. the lines through the scaled points $x^*(a)$ and $a$, where $a$ is chosen arbitrarily (so that $na$ is outside of the scaled domain $D$).

In order to justify the tangent method, we need the following.

**Theorem 2.1.** Let us consider the set $P_{a,b}$ of weighted paths from the origin $(0,0)$ to a point $(a,b)$ on the $\mathbb{Z}^2$ lattice, that use a finite family of steps $\{(s_i, t_i), i = 1, 2, \ldots, k\}$ with $t_i > 0$ for all $i$ (directed paths moving to the right), with respective weights $w_i, i = 1, 2, \ldots, k$. In this model each path $p \in P_{a,b}$ is weighted by $w(p) = \prod_{s \in p} w(s)$, where $w(s) = w_i$ if $s = (s_i, t_i)$. In other words, defining the partition function of the model as

$$Z_{(0,0) \to (a,b)} = \sum_{p \in P_{a,b}} w(p)$$

we consider the paths of $P_{a,b}$ as a statistical ensemble with probability weights $P(p) = w(p)/Z_{(0,0) \to (a,b)}$. Then the most likely configuration of a weighted path in $P_{n^\alpha, n^\beta}$ for large $n$ is the straight line from the origin to $(a,b) = n(\alpha, \beta)$.

**Proof.** Let $P$ denote the Newton polynomial of the collection of weighted steps:

$$P(x, y) := \sum_{i=1}^k w_i x^{s_i} y^{t_i}.$$ 

Then the generating partition function for weighted paths from $(0,0)$ to $(a,b)$, with an additional weight $t$ per step reads:
\[ f_{a,b}(t) = \sum_{\text{paths } p \in P_{a,b}} w(p) f(p) \frac{1}{1 - t P(x,y)} \bigg|_{x=\beta} = \int \frac{dx dy}{(2\pi)^2 x^{a+1} y^{b+1}} \frac{1}{1 - t P(x,y)} \]

where \( f(p) \) is the length of \( p \), i.e. its total number of steps, and where we use the notation \( f(x,y) \big|_{x=\beta} \) for the coefficient of \( x^a y^b \) in the series \( f(x,y) \) (here we expand the fraction as a series of \( t \)). We may split the corresponding partition function \( Z_{(0,0) \to (a,b)} \) into two pieces by recording the position of some intermediate point \((x,y)\) in a domain \( L \), resulting in the decomposition:

\[ Z_{(0,0) \to (a,b)} = \sum_{(x,y) \in L} Z_{(0,0) \to (x,y)} Z_{(x,y) \to (a,b)}. \]

Let us now consider the scaling limit where \((a,b) = n(\alpha, \beta)\), and \((x,y) = n(\xi, \eta)\), \( L = nL \), with very large \( n \). We have asymptotically

\[ Z_{(0,0) \to n(\alpha, \beta)} \sim \int_L d\xi d\eta \ Z_{(0,0) \to n(\xi, \eta)} Z_{n(\xi, \eta) \to n(\alpha, \beta)} \]

where

\[ Z_{(0,0) \to n(\alpha, \beta)} = \int \frac{dx dy}{(2\pi)^2 x^a y^b} e^{nS_{\alpha, \beta}(x,y)}, \quad S_{\alpha, \beta}(x,y) = -\alpha \log(x) - \beta \log(y) - \frac{1}{n} \log(1 - t P(x,y)) \]

and

\[ Z_{n(\xi, \eta) \to n(\alpha, \beta)} = Z_{(0,0) \to n(\alpha - \xi, \beta - \eta)} \]

by translational invariance of the problem. We may rewrite

\[ Z_{(0,0) \to n(\alpha, \beta)} \sim \int \frac{dx dy}{(2\pi)^2 x^a y^b} e^{nS}, \quad S = S_{\xi, \eta}(x,y) + S_{\alpha - \xi, \beta - \eta}(u,v). \]

The integral is dominated by the extrema of the action. Writing \( \partial_\xi S = \partial_\eta S = 0 \) implies \( x = u \) and \( y = v \). Moreover we have:

\[
\begin{align*}
\partial_\xi S &= \frac{\xi}{x} + \frac{1}{n} \frac{t \partial_x P(x,y)}{1 - t P(x,y)} = 0 \\
\partial_\eta S &= -\frac{\eta}{y} + \frac{1}{n} \frac{t \partial_y P(x,y)}{1 - t P(x,y)} = 0 \\
\partial_u S &= \frac{\alpha - \xi}{u} + \frac{1}{n} \frac{t \partial_u P(u,v)}{1 - t P(u,v)} = 0 \\
\partial_v S &= -\frac{\beta - \eta}{v} + \frac{1}{n} \frac{t \partial_v P(u,v)}{1 - t P(u,v)} = 0.
\end{align*}
\]

At the saddle-point, as \( u = x \) and \( v = y \), we deduce that:

\[
\frac{\xi}{\eta} = \frac{\alpha - \xi}{\beta - \eta} \Rightarrow \beta \xi - \alpha \eta = 0.
\]
We conclude that the most likely intermediate point lies on the line through $(0,0)$ and $(\alpha, \beta)$, and the theorem follows.

To make the tangent method completely rigorous, we should argue that the path is most likely to follow a straight line until the contact with the arctic curve, when the influence of the other paths starts to play a role. One also should argue that the reasoning holds irrespectively of possibly stronger interactions between paths (like in the six vertex case).

2.2. Non-intersecting paths and LU method

Most of the models that we consider in this paper, with the exception of the vertically symmetric alternating sign matrices in section 7, have a formulation in terms of NILPs on a directed graph $\mathcal{G}$. The general setup is as follows. To start, we may allow the paths to intersect, and we only impose the non-intersecting condition after setting down the general definitions. We consider $n+1$ lattice paths on $\mathcal{G}$, with oriented steps along the oriented edges of $\mathcal{G}$, which begin at the vertices $v_i$ and end at vertices $\tilde{v}_j$, with $i,j$ integers in the range $0, \ldots, n$. For each oriented edge $e$ of the graph we assign a weight $w(e)$. To each individual path $p$ we assign the weight $w(p) = \prod_{e \in p} w(e)$.

For a collection of $n+1$ lattice paths $P = \{p_0, p_1, \ldots, p_n\}$ we assign the weight $w(P) = \prod_{j=0}^{n} w(p_j)$. Finally, we define the partition function for lattice paths on the graph $\mathcal{G}$ which start at the vertices $\{v_i\}$ and end at $\{\tilde{v}_j\}$ to be

$$Z(\{v_i\} \to \{\tilde{v}_j\}) = \sum_{P: \{v_i\} \to \{\tilde{v}_j\}} w(P).$$

The sum is taken over all sets $P$ of $n+1$ lattice paths on $\mathcal{G}$ which connect the starting vertices $v_i$ to the ending vertices $\tilde{v}_j$.

To address the case of non-intersecting lattice paths, we shall use the celebrated Gessel–Viennot formula [8]. However it is only applicable if our path setting satisfies the following crossing property:

$$(CP) \text{ For any } 0 \leq i < j \leq n \text{ and } 0 \leq k < \ell \leq n, \text{ any path from } v_i \text{ to } \tilde{v}_\ell \text{ and any path from } v_j \text{ to } \tilde{v}_k \text{ must intersect at least once.}$$

This puts in principle some restriction on the choice of starting and ending points, depending on the structure of the directed graph $\mathcal{G}$. However for all the applications in the present paper, the crossing property will always be trivially satisfied.

**Lemma 2.2 ([8]).** For non-intersecting lattice paths satisfying the property (CP), the partition function $Z(\{v_i\} \to \{\tilde{v}_j\})$ is given by the determinant formula

$$Z(\{v_i\} \to \{\tilde{v}_j\}) = \det_{i,j \in [0,n]} (Z_{ij}). \quad (2.2)$$

where

$$Z_{ij} = \sum_{p: v_i \to \tilde{v}_j} w(p). \quad (2.3)$$

is the weight for all paths from $v_i$ to $\tilde{v}_j$.

In applying the tangent method to non-intersecting path models we encounter the following situation. We start with the original model of non-intersecting paths on a certain domain $D$ of the directed lattice graph $\mathcal{G}$. Its partition function $Z$ is given by lemma 2.2 as the determinant of a certain matrix $A$, $Z = \det(A)$. 


Next, we must consider the partition function \( N_\tilde{v} \) of the same model, but with the endpoint \( \tilde{v}_n \) moved to a different location \( \tilde{v} \) outside of \( D \), such that the Gessel–Viennot formula of lemma 2.2 still applies. This partition function may be decomposed according to the position \( \tilde{w} \) (on the boundary \( \partial D \) of \( D \)) of the exit point from the domain \( D \) of the path from \( v_n \to \tilde{v} \) (we choose the point \( \tilde{v} \) so that the path only exits once, and can never visit \( D \) again). Let \( Z_\tilde{w} \) denote the partition function of paths on \( D \) from \( \{ v_0, v_1, ..., v_n \} \) to \( \{ \tilde{v}_0, \tilde{v}_1, ..., \tilde{v}_{n-1}, \tilde{w} \} \), and \( Y_{\tilde{w}, \tilde{v}} \) that for paths from \( \tilde{w} \) to \( \tilde{v} \) that exit \( D \). Then we have

\[
N_\tilde{v} = \sum_{\tilde{w} \in \partial D} Z_\tilde{w} Y_{\tilde{w}, \tilde{v}}. \tag{2.4}
\]

Applying the Gessel–Viennot formula of lemma 2.2, we find that \( Z_\tilde{w} \) is the determinant of a matrix \( A^{(\tilde{w})} \) which differs from \( A \) only in its last column, \( Z_\tilde{w} = \det(A^{(\tilde{w})}) \). For each \( \tilde{w} \) we have

\[
A_{ij}^{(\tilde{w})} = \begin{cases} A_{ij}, & j \in [0, n-1] \\ b_i^{(\tilde{w})}, & j = n \end{cases}, \tag{2.5}
\]

where \( A_{ij} = Z_{v_i \to v_j} \) and \( b_i^{(\tilde{w})} = Z_{v_i \to \tilde{w}} \) is the partition function for the paths from \( v_i \) to \( \tilde{w} \) on \( D \).

To compute the determinant of \( A^{(\tilde{w})} \) we use the LU decomposition. First, let us suppose that the original matrix \( A \) has an LU decomposition as

\[
A = LU, \tag{2.6}
\]

where \( L \) is a lower triangular matrix with 1’s on the diagonal and \( U \) is an upper triangular matrix. Then we have

\[
\det_{i,j \in [0,n]} (A_{ij}) = \prod_{i=0}^{n} U_{ii}. \tag{2.7}
\]

In the cases we consider in this paper the LU decomposition of the matrix \( A \) is either known and can be found in the literature, or we compute the matrices \( L \) and \( U \) explicitly using various methods. This gives immediately an LU decomposition for the matrices \( A^{(\tilde{w})} \) as well:

\[
A^{(\tilde{w})} = LU^{(\tilde{w})}, \tag{2.8}
\]

where

\[
U^{(\tilde{w})} := L^{-1} A^{(\tilde{w})}. \tag{2.9}
\]

It then follows that

\[
\det_{i,j \in [0,n]} (A_{ij}^{(\tilde{w})}) = \prod_{i=0}^{n} U_{ij}^{(\tilde{w})}. \tag{2.10}
\]

The advantage of this method is that a major simplification occurs due to the fact that \( A^{(\tilde{w})} \) differs from \( A \) only in the last column. Indeed, we have \( U_{ij}^{(\tilde{w})} = \sum_{k=0}^{n} (L^{-1})_{ik} A_{kj}^{(\tilde{w})} \), but \( A_{kj}^{(\tilde{w})} = A_{kj} \) for \( j < n \). This means that for \( j < n \), we have \( U_{ij}^{(\tilde{w})} = U_{ij} \). Thus, we find that the determinant of \( A^{(\tilde{w})} \) is given by

\[
\det_{i,j \in [0,n]} (A_{ij}^{(\tilde{w})}) = \left( \prod_{i=0}^{n-1} U_{ii} \right) U_{nn}^{(\tilde{w})}. \tag{2.11}
\]

In other words, the computation of the determinant of \( A^{(\tilde{w})} \) reduces to the computation of the single matrix element
\[ U^{(\tilde{w})}_{n,m} = \sum_{k=0}^{n} (L^{-1})_{n,k} A^{(\tilde{w})}_{k,n} = \sum_{k=0}^{n} (L^{-1})_{n,k} b^k_{(\tilde{w})} . \]  

(2.12)

In the context of the tangent method it is more useful to consider a ‘one-point function’ which is defined as the following ratio \( H_{\tilde{w}} \) of partition functions:

\[ H_{\tilde{w}} := \frac{Z_{\tilde{w}}}{Z} = \frac{\det(A^{(\tilde{w})})}{\det(A)} . \]  

(2.13)

Using the LU decomposition method we find that this function is given by

\[ H_{\tilde{w}} = \frac{U^{(\tilde{w})}_{n,n}}{U^{(\tilde{w})}_{n,n}} . \]  

(2.14)

This gives an efficient way for calculating \( H_{\tilde{w}} \). The tangent method can then be applied to the decomposition of the normalized partition function, in the limit of a large scaled domain \( D \):

\[ \frac{N_{\tilde{w}}}{Z} = \sum_{\tilde{w}} H_{\tilde{w}} Y_{\tilde{w},\tilde{\tau}} . \]  

(2.15)

However, to perform an asymptotic analysis of this sum, we need to be able to estimate \( H_{\tilde{w}} \) in large size. Note that the expression for \( U^{(\tilde{w})}_{n,n} \) obtained from the computation of \( L^{-1} A^{(\tilde{w})} \) is naturally an alternating sum due to the cofactor expansion of \( L^{-1} \). Such sums are not suitable for asymptotic analysis, as large terms may cancel out. We will therefore need to reexpress \( H_{\tilde{w}} \) as positive sums, a technically demanding process.

2.3. Working with generating functions

In the core of the paper, we will use a number of manipulations involving generating functions which we gather here as a preliminary.

2.3.1 Generating functions and binomial identities. For a series or polynomial \( f(x) = \sum_{n \geq 0} f_n x^n \) we write \( f_n = f(x)|_{x^n} \) for the coefficient of \( x^n \) in \( f(x) \) (in particular when \( n = 0 \) this denotes the constant term \( f_0 \) of \( f \)). We also have by the Cauchy theorem:

\[ f(x)|_{x^n} = \oint \frac{dx}{2\pi i x^{n+1}} f(x) \]  

(2.16)

where the contour of integration is around 0, and the integral picks up the residue at 0.

The following lemma will be used throughout the paper.

**Lemma 2.3.** For any two generating functions \( f(x), g(x) \), we have the following identity:

\[ f(x)g(x)|_{x^\epsilon} = \sum_{\ell=0}^k f(x)|_{x^\ell} g(x)|_{x^{\epsilon-\ell}} . \]  

(2.17)

**Proof.** This expresses simply \( (fg)_n \), the coefficient of \( x^n \) in \( fg \), as \( (fg)_n = \sum_{\ell=0}^k f_n g_{k-n} \).}

In this paper, we will deal with expressions involving typically sums of products of binomial coefficients. The following lemma will be used repeatedly.
**Lemma 2.4.** We have four simple ways of expressing the binomial coefficient \( \binom{n}{k} \) as the coefficient in a generating series or polynomial:

\[
\binom{n}{k} = \left. (1 + x)^n \right|_{x^k} = \left. \frac{1}{(1 - x)^{k+1}} \right|_{x^{n-k}} = \frac{1}{(1 - x)^{n-k+1}}.
\]

**(Proof.***) The two first expressions are obvious. To get the third and fourth, simply notice that

\[
\binom{n}{k} = \left. (1 - x)^m \right|_{x^n} = \binom{m + j}{j} x^j.
\]

\[(2.18)\]

We shall also use the description below, using iterated derivatives.

**Lemma 2.5.** We have the following representations of the binomial coefficient \( \binom{n}{k} \):

\[
\binom{n}{k} = \frac{1}{k!} \left. \frac{d^k}{dx^k} r^p \right|_{r=1} = \frac{1}{(n - k)!} \left. \frac{d^{n-k}}{dx^{n-k}} r^p \right|_{r=1}.
\]

\[(2.20)\]

The collection of lemmas above will be used extensively throughout this paper to evaluate summations over products of binomial coefficients.

A last formula concerns the description of the inverse of a binomial coefficient, by means of a hypergeometric series. By these we mean precisely the series

\[
_2F_1(a, b, c; x) := \sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{(c)_n n!} x^n = \int_{0}^{1} dr \frac{r^{b-1}(1 - r)^{c-b-1}(1 - tx)^{-a}}{1 - tx}
\]

where we used the notation \((x)_n = x(x-1)...(x-n+1)\) for the Pochhammer symbol. We have the following direct application:

**Lemma 2.6.** The inverse of the binomial coefficient \( \binom{n+a}{a} \) for \( n, a \geq 0 \) is given by:

\[
\binom{n+a}{a}^{-1} = a \times _2F_1(1, 1, a + 1; x)\big|_{x=1}.
\]

Using the explicit integral formulation above, we may also write:

\[
\sum_{n=0}^\infty \binom{n+a}{a}^{-1} x^n = a \int_{0}^{1} dr \frac{(1 - t)^{a-1}}{1 - tx} = a \int_{0}^{1} dt \frac{t^{a-1}}{1 - x + tx} = a \sum_{m=0}^\infty (-1)^m \frac{x^m}{(1 - x)^{m+1}} \int_{0}^{1} dr r^{m+a-1} = \sum_{m=0}^\infty \frac{(-1)^m a}{m + a} \frac{x^m}{(1 - x)^{m+1}}.
\]

Extracting the coefficient of \( x^n \) we get the following formula:

**Lemma 2.7.** The inverse of the binomial coefficient \( \binom{n+a}{a} \) for \( n, a \geq 0 \) is given by:

\[
\binom{n+a}{a}^{-1} = \sum_{m=0}^{n} \frac{(-1)^m a}{m + a} \binom{n}{m}.
\]
2.3.2. Infinite matrices and their truncations. In this paper, we will also deal with some finite truncations of infinite matrices, whose manipulation is greatly simplified by use of generating functions. For an infinite matrix \( A \) with entries \( A_{ij}, i, j \in \mathbb{Z}_{\geq 0} \), we introduce the generating function

\[
f_A(z, w) := \sum_{ij \geq 0} z^i w^j A_{ij}.
\]

Then we have the following result.

**Lemma 2.8.** For any two infinite matrices \( A, B \) with generating functions \( f_A, f_B \), assuming the product \( AB \) makes sense, we have the following product formula:

\[
f_{AB}(z, w) = f_A \ast f_B(z, w) = f_A(z, t^{-1}) f_B(t, w)|_t = \oint \frac{dt}{2i\pi} f_A(z, t^{-1}) f_B(t, w)
\]

where \( \ast \) stands for the convolution product of two-variable generating functions, and the contour integral is for instance over the unit circle.

**Proof.** We write:

\[
f_{AB}(z, w) = \sum_{ij \geq 0} z^i w^j \sum_{r \geq 0} A_{ir} B_{rj} = f_A(z, t^{-1}) f_B(t, w)|_t.
\]

Another important property of infinite matrices is that the LU factorization of an infinite matrix \( A \) descends to its finite truncations \( A^{(n)} = (A_{ij})_{ij \in [0,n]} \). More precisely, we have the following:

**Lemma 2.9.** Let \( A \) be an infinite matrix. Assume there exist infinite matrices \( L, U \) respectively uni-lower- and upper-triangular such that \( A = LU \). Then for all \( n \geq 0 \) we have the following LU factorization:

\[
A^{(n)} = L^{(n)} U^{(n)} \tag{2.21}
\]

and moreover if \( L \) is invertible, its inverse also truncates:

\[
(L^{(n)})^{-1} = (L^{-1})^{(n)}. \tag{2.22}
\]

**Proof.** Assuming \( A = LU \), we compute:

\[
A^{(n)} = \left( \sum_{r \geq 0} L_{ir} U_{rj} \right)_{ij \in [0,n]} = \left( \sum_{r=j} L_{ir} U_{rj} \right)_{ij \in [0,n]} = \left( \sum_{r=j} L^{(n)}_{ir} U^{(n)}_{rj} \right)_{ij \in [0,n]}
\]

where the truncation appears de facto from triangularity and (2.21) follows. Assuming \( L^{-1} \) exists, it is also uni-triangular, and we write similarly:

\[
I^{(n)} = (L^{-1} L)^{(n)} = \left( \sum_{r \geq 0} (L^{-1})_{ir} L_{rj} \right)_{ij \in [0,n]} = \left( \sum_{r \geq 0} (L^{-1})_{ir} L^{(n)}_{rj} \right)_{ij \in [0,n]}
\]

\[
= \left( \sum_{r \geq 0} \left( (L^{-1})^{(n)}_{ir} \right) L^{(n)}_{rj} \right)_{ij \in [0,n]}
\]

which implies (2.22).
3. Aztec diamond domino tilings

3.1. Path formulation

The Aztec diamond domino tiling problem can be stated as follows: tile the size \( n \) domain \( D \) of the square lattice \( \mathbb{Z}^2 \) depicted in figure 2 by means of dominos made of two elementary squares glued along a common edge. After bi-coloring the square faces of \( D \) (in black and white), so that South–West (SW) boundary squares are colored in black, we see that there are 4 distinct bi-colored tiles compatible with the bi-coloration of the faces of \( D \).

Domino tiling configurations of \( D \) are in bijection with non-intersecting paths. Simply associate a portion of path drawn on the four possible domino tiles:

\[
\text{(3.1)}
\]

These paths join vertices at the middle of edges of the original square lattice, and are non-intersecting by definition. Note that some regions may not be visited by paths, due to the fact that the third domino above is not crossed by any path. For the Aztec diamond of size \( n \), the domino tiling configurations are equivalent to those of \( n+1 \) NILPs with steps \((1, 1)\), \((1, -1)\) and \((2, 0)\) which we call respectively up, down and horizontal, corresponding to the 1st, 2nd and 3rd dominos above. Moreover \( n \) of the paths start at the middle of the West edge of the SW boundary black square faces introduced above. We may choose coordinates in which these are \((-i, i), i = 1, 2, ..., n\). These paths end on the SE boundary white square faces at points with coordinates \((j, j), j = 1, 2, ..., n\). In addition, we consider the origin to be both the starting and end point of a trivial path of length 0. The non-intersection constraint then forces any path from \((-1, 1)\) to \((1, 1)\) to have either a horizontal step or a succession up-down, but prevents the unwanted down-up succession. Such paths are usually referred to as large Schroeder paths. If we assign a weight of 1 to each edge of the lattice, then the partition function of this system is equal to the number \(2^{n(n+1)/2}\) of domino tilings of the Aztec diamond of size \( n \). We now briefly review this result.

3.2. Exact enumeration

We now review the calculation, in the paths formulation, of the number of domino tilings of the Aztec diamond. First, if we assign a weight of 1 to each edge of the lattice, then the partition function for all paths from \((-i, i)\) to \((j, j)\) is readily calculated to be

\[
Z_{(-i, i)\to(j, j)} = A_{ij} := \sum_{p=0}^{\min(i,j)} \binom{i+j-p}{p, i-p, j-p}
\]

(3.2)

where in the summation \( p \) denotes the number of horizontal steps, and \( i-p, j-p \) the numbers of up and down steps, respectively. Then by lemma 2.2, the partition function for this system is \( Z = \det(A) \). This determinant can be computed with the help of the following lemma.
Lemma 3.1. The matrix $A$ admits an LU decomposition $A = LU$ with
\[
L_{ij} = \binom{i}{j},
\]
and
\[
U_{ij} = 2^i \binom{i}{j}.
\]

Since $L_{jj} = 1$ and $U_{jj} = 2^j$, an immediate consequence of the lemma is that
\[
\det \left( A_{ij} \right) = \prod_{i=0}^{n} U_{ii} = 2^{\frac{n(n+1)}{2}},
\]
which is exactly equal to the number of domino tilings of the Aztec diamond of size $n$.

Proof. To prove lemma 3.1 we use the infinite matrix generating function methods of section 2.3.2. More precisely, we find a simple LU decomposition of the infinite matrix $A = (A_{ij})_{i,j \in \mathbb{Z}_{\geq 0}}$, with generating function:
\[
f_A(z,w) := \sum_{i,j=0}^{\infty} A_{ij} z^i w^j = \frac{1}{1 - z - w - zw}.
\]

This is proved by working backwards from the generating function. We have
\[
\frac{1}{1 - z - w - zw} = \sum_{m=0}^{\infty} (z + w + wz)^m \\
= \sum_{m=0}^{\infty} \sum_{p_1, p_2, p_3 = 0}^{\infty} \left( \begin{array}{c} m \\ p_1, p_2, p_3 \end{array} \right) z^{p_1} w^{p_2} (zw)^{p_3} \\
= \sum_{p_1, p_2, p_3 = 0}^{\infty} \left( p_1 + p_2 + p_3 \right) z^{p_1+p_2+p_3},
\]

by the standard trinomial identity. To extract the coefficient of \(z^i w^j\) in this we set \(p_1 = i - p_3\), \(p_2 = j - p_3\) to find:

\[
\frac{1}{1 - z - w - zw} \bigg|_{z^i w^j} = \sum_{p_3 = 0}^{\min(i,j)} \left( \begin{array}{c} i + j - p_3 \\ i - p_3, j - p_3, p_3 \end{array} \right),
\]

(3.7)

which coincides with our original expression for the weights \(A_{i,j}\).

A particular LU decomposition of \(A\) is obtained by checking that:

\[
f_A(z, w) = (f_L \star f_U)(z, w) \tag{3.8}
\]

where

\[
f_L(z, w) = \frac{1}{1 - z - zw} \quad \text{and} \quad f_U(z, w) = \frac{1}{1 - w - 2zw}. \tag{3.9}
\]

Indeed, we have:

\[
(f_L \star f_U)(z, w) = \oint \frac{dr}{2i\pi r (1 - z)} \frac{1}{1 - w - 2zw} = \frac{1}{1 - z} \frac{1}{1 - w - 2zw} = \frac{1}{1 - z - w - zw}
\]

where we noted that the contour integral picks up the residue at \(t = z/(1 - z)\). Comparing this with (3.6) implies that the matrix \(A\) factorizes as \(A = LU\) where the \((i,j)\) matrix elements of \(L\) and \(U\) are given by the coefficient of \(z^i w^j\) in the expansion of the generating functions \(f_L(z, w)\) and \(f_U(z, w)\) about \(z = w = 0\). Explicitly, we have

\[
f_L(z, w) = \sum_{i=0}^{\infty} (1 + w)^i z^i = \sum_{i=0}^{\infty} \sum_{j=0}^{i} \binom{i}{j} z^i w^j, \tag{3.10}
\]

from which we find that \(L\) is lower uni-triangular, with entries \(L_{ij} = \binom{i}{j}\). Similarly, we have:

\[
f_U(z, w) = \sum_{j=0}^{\infty} (1 + 2z)^j w^j = \sum_{j=0}^{\infty} \sum_{i=0}^{j} 2^j \binom{j}{i} z^i w^j \tag{3.11}
\]

from which we get that \(U\) is upper triangular with entries \(U_{ij} = 2^j \binom{j}{i}\). As noted in section 2.3.2, this \(LU\) decomposition holds for the finite matrices obtained by truncation to indices \(i, j \in [0, n]\), and lemma 3.1 follows (Note that we dropped the superscript \((n)\) in \(A, L, U\) for simplicity).
For later use we note that the infinite matrix $L$ is invertible, and the inverse matrix $L^{-1}$ has matrix elements

$$(L^{-1})_{i,j} = (-1)^{i+j} \binom{i}{j}$$

(3.12)

and generating function

$$f_{L^{-1}}(z,w) = \frac{1}{1+z-zw}.$$  

(3.13)

Indeed, we easily compute:

$$(f_{L^{-1}} * f_L)(z,w) = \oint_{2\pi i} \frac{dt}{2\pi i} \frac{1}{t(1+z) - z} \frac{1}{1-t(1+w)} = \frac{1}{1+z} \frac{1}{1-z} = \frac{1}{1-zw} = f_L(z,w)$$

where the contour integral has picked the residue at $t = z/(1+z)$. By lemma 2.9, equation (3.12) also holds for the finite truncation of $L$ to indices $i,j \in [0,n]$.

As explained in section 2.1, to set up the tangent method for the Aztec diamond we extend the original diamond shaped domain by appending a rectangular region with corners at $(n,n)$, $(0,2n)$, $(k-n,k+n)$ and $(k,k)$, for some integer $k > n$. We again consider $n + 1$ paths on this domain with the same starting points $(-i,i)$, $i \in [0,n]$ as before, but now we take the end points to be $(j,j)$, $j \in [0,n-1]$, and $(k,k)$. Thus, in the non-intersecting configuration the top-most path will start at $(-n,n)$ and end at $(k,k)$ at the far upper right corner of the extended domain.

Let us denote the partition function for these $n + 1$ paths on the extended domain by $N(n,k)$. To apply the tangent method we expand $N(n,k)$ in terms of position $(\ell,2n-\ell)$, $\ell \in [0,n]$, of the point where the top-most path exits from the original domain $D$ into the extended domain.

We have

$$N(n,k) = \sum_{\ell=0}^{n} Z(n,\ell) Y(\ell,k),$$

(3.14)

where $Z(n,\ell)$ is the partition function for $n + 1$ paths on the original domain with starting points at $(-i,i)$, $i \in [0,n]$ and ending points at $(-j,j)$, $j \in [0,n-1]$ and $(\ell,2n-\ell)$, and $Y(\ell,k)$ is the partition function for the single path exiting $D$ at $(\ell,2n-\ell)$, and ending at $(k,k)$. In computing $Y(\ell,k)$ we should only include contributions from paths which start with a step of the form $(1,1)$, or $(2,0)$ (otherwise the path would not leave $D$). It is therefore the sum of two contributions, one for paths $(\ell + 1,2n - \ell + 1) \to (k,k)$ and one for paths $(\ell + 2,2n - \ell) \to (k,k)$. By translation invariance the paths from $(a,b)$ to $(c,d)$ contribute the same as those from $(0,0) \to (c-a,d-b)$, with partition function $A_{ij}$ (3.2), where $i$ and $j$ are such that $i+j = c-a$ and $j-i = d-b$, hence $i = \frac{b+c-a-d}{2}$ and $j = \frac{c+d-a-b}{2}$. As a consequence, $Y(\ell,k)$ is given by:

$$Y(\ell,k) = A_{n-\ell,k-n-1} + A_{n-\ell-1,k-n-1}.$$  

(3.15)

The partition function $Z(n,\ell)$ is given by lemma 2.2 as the determinant of a certain $(n + 1) \times (n + 1)$ matrix $A$ which is constructed as follows. First note that the weight for all paths from $(-i,i)$ to $(\ell,2n-\ell)$ is by translational invariance the same as that from $(0,0) \to (\ell+i,2n-\ell-i)$, hence:
Using generating functions, we may write:

\[
\tilde{U}_{n,n} = \sum_{r=0}^{n} (L^{-1})_{n,r} \tilde{A}_{r,n} = \sum_{r=0}^{n} (-1)^{n+r} \binom{n}{r} A_{r+n-\ell,n} = \sum_{r=0}^{n} (-1)^{n-r} \binom{n}{r} \frac{1}{1-z-w-2w} \left| L^{-1}_{r+n-\ell+\ell-\ell,n} \right|.
\]

Theorem 3.2. The one point function \( H(n, \ell) \), 0 \( \leq \ell \leq n \), for the Aztec diamond tiling problem is

\[
H(n, \ell) = \frac{1}{2^\ell} \sum_{p=0}^{\ell} \binom{n}{p}.
\]

Proof. From section 2.2 we have \( H(n, \ell) = \frac{U_{n,n}}{U_{\ell,n}} \), where \( \tilde{U} := L^{-1} \tilde{A} \) and \( L^{-1} \) was given in equation (3.12). In addition, \( U_{n,n} = 2^n \), so we only need to prove that:

\[
\tilde{U}_{n,n} = \sum_{p=0}^{\ell} \binom{n}{p}.
\]

Using generating functions, we may write:

\[
\tilde{U}_{n,n} = \sum_{r=0}^{n} (L^{-1})_{n,r} \tilde{A}_{r,n} = \sum_{r=0}^{n} (-1)^{n+r} \binom{n}{r} A_{r+n-\ell,n} = \sum_{r=0}^{n} (-1)^{n-r} \binom{n}{r} \frac{1}{1-z-w-2w} \left| L^{-1}_{n-\ell+n-\ell,n} \right|.
\]

and the theorem follows.
3.3. Tangent method

Following section 2.1, we wish to evaluate the function $N(n,k)$ (3.14) in the large size limit corresponding to the scaling $k = nz$, $n$ large and $z > 1$ independent of $n$. The leading large $n$ asymptotics are governed by:

$$\frac{N(n,nz)}{Z(n,n)} \sim \int_0^1 d\xi \; H(n,n\xi) Y(n\xi,nz),$$

where we have set $\ell = n\xi$ and replaced the summation by an integral. From the explicit formulas (3.15) and (3.19), and using the Stirling formula, we get the following asymptotic behaviors:

$$Y(n\xi, nz) \sim 2A_n(1-\xi)n(z-1) \sim \int_0^{\Min(1-\xi,z-1)} d\theta \; e^{nS_0(\xi,\theta; z)},$$

where

$$S_0(\xi,\theta; z) = (z - \xi - \theta)\Log(z - \xi - \theta) - \theta\Log(\theta) - (1 - \xi - \theta)\Log(1 - \xi - \theta) - (z - 1 - \theta)\Log(z - 1 - \theta)$$

and

$$H(n, n\xi) \sim \frac{1}{2n} \int_0^\xi d\varphi \; e^{nS_1(\varphi; z)},$$

where

$$S_1(\varphi; z) = -\varphi\Log(\varphi) - (1 - \varphi)\Log(1 - \varphi).$$

The extremum of the latter action is at $\varphi = \varphi^* = \frac{1}{2}$. Two cases must be distinguished:

- (i) $\xi > \frac{1}{2}$: then the asymptotics of $H(n, n\xi)$ are dominated by $e^{nS_1(\varphi^*; z)/2^n} \sim 1$ hence no contribution to the large $n$ dominant behavior.
- (ii) $\xi < \frac{1}{2}$: the integral is dominated by the value at its upper bound $\xi$, namely

$$H(n, n\xi) \sim e^{nS_1(\xi; z) - n\Log(2)}.$$

Collecting both results, we find that $\frac{N(n,nz)}{Z(n,n)}$ is dominated by $e^{nS_0(\xi,\theta; z)}$, where:

- (i) If $\xi > \frac{1}{2}$: $S(\xi,\theta; z) = S_0(\xi,\theta; z)$.
- (ii) If $\xi < \frac{1}{2}$: $S(\xi,\theta; z) = S_0(\xi,\theta; z) + S_1(\xi; z) - \Log(2)$.

Note that

$$\partial_\xi S_0(\xi,\theta; z) = \Log\left(\frac{(1 - \xi - \theta)(z - 1 - \theta)}{\theta(z - \xi - \theta)}\right) \quad \text{and} \quad \partial_\theta S_0(\xi,\theta; z) = \Log\left(\frac{1 - \xi - \theta}{z - \xi}\right).$$

The extremization problem in the case (i), $\partial_\theta S_0(\xi,\theta; z) = \partial_\xi S_0(\xi,\theta; z) = 0$ has clearly no solution as one should have $\theta = 1 - z$, but $z > 1$. We are left with (ii). Writing $\partial_\xi S = \partial_\theta S = 0$, we get:

$$(1 - \xi - \theta)(z - 1 - \theta) = \theta(z - \xi - \theta) \quad \text{and} \quad (1 - \xi - \theta)(1 - \xi) = (z - \xi - \theta)\xi.$$ 

Eliminating $\theta$ yields $\xi = \xi^*(z) := \frac{1}{2z} < \frac{1}{2}$ as $z > 1$. 

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The line through the scaled points \((\xi^*(z), 2 - \xi^*(z))\) and \((z, z)\) has the equation:

\[
y = \frac{2 - \xi^*(z) - z}{\xi^*(z) - z} x - 2z \frac{1 - \xi^*(z)}{\xi^*(z) - z}.
\] (3.21)

Following section 2.1, the arctic curve must be the envelope of this family of lines, hence it is determined by the equation (3.21) together with its derivative w.r.t. \(z\). This gives the system:

\[
\begin{align*}
-x - y + 2z + 4xz - 4z^2 - 2xz^2 + 2yz^2 &= 0 \\
1 + 2x - y - 2xy - 3z - 2xz + 2x^2z + 2yz + 2xyz &= 0
\end{align*}
\]

whose solution is given the parametric equations, for \(z > 1\):

\[
\begin{align*}
x &= \frac{1}{2} - \frac{(z - 1)(2z - 1)}{2z(z - 1) + 1}, \\
y &= \frac{3}{2} + \frac{z - 1}{2z(z - 1) + 1}.
\end{align*}
\]

Alternatively, eliminating \(z\), we end up with the arctic circle:

\[
x^2 + (y - 1)^2 = \frac{1}{2}
\] (3.22)

centered at \((0, 1)\) with radius \(1/\sqrt{2}\), namely inscribed in the limiting scaled domain \(\lim_{n \to \infty} D/n\) i.e. the square defined by the equation \(|x| + |y - 1| = 1\). *Stricto sensu*, we have only derived the portion of the arctic circle corresponding to \(z > 1\), namely \(-\frac{1}{2} < x < \frac{1}{2}\) and \(\frac{3}{2} < y < 2\). However, due to the symmetry of the problem under rotation of \(\pi/2\), the same analysis can be performed on the rotated domains, leading to (3.22) for all \(x, y\).

### 4. The non-intersecting Dyck path problem

In this section we study a different problem of non-intersecting lattice paths, and apply the tangent method to determine the corresponding arctic curve.

#### 4.1. Counting problem

We consider the situation of figure 3. We wish to count the number \(N(n, k, p)\) of path configurations with up/down steps \((1, \pm 1)\) only, that remain above the \(x\) axis, and start at points \((-2i, 0)\), \(i = 0, 1, .., n\), and end at points \((2k + 2i, 0)\), \(i = 0, 1, .., n - 1\) and \((2k + n + p, n + p)\). Denote by \(Z(n; k; \ell)\) the partition function corresponding to the lower \((n + k) \times (n + k)\) truncated square, and \(Y(p; \ell)\) that of the upper \(p \times (n + k)\). Then we have

\[
N(n, k, p) = \sum_{\ell=0}^{n+k} Z(n, k; \ell) Y(p; \ell).
\] (4.1)

We shall rather consider the ‘one-point function’

\[
H(n, k; \ell) := \frac{Z(n, k; \ell)}{Z(n, k; 0)},
\] (4.2)

and we will apply the saddle-point method to the sum:

\[
\frac{N(n, k, p)}{Z(n, k; 0)} = \sum_{\ell=0}^{n+k} H(n, k; \ell) Y(p; \ell).
\] (4.3)
The latter partition function $Y(p; \ell)$ is very simple, as it enumerates the configurations of a single path from $(2k + n + 1 - \ell, n + 1 + \ell)$ to $(2k + n + p, n + p)$, namely

$$Y(p; \ell) = \binom{p + \ell - 1}{\ell}.$$  \hfill (4.4)

The next two sections are devoted to the computation of $H(n, k; \ell)$.

### 4.2. Partition function

In this section, we compute the partition function $Z(n, k; 0)$. Note that it may be obtained by specializing a result from 1989 by Gessel and Viennot (see Krattenthaler [15] for details and proofs), however the method of proof is completely different. Here we derive the result using LU decomposition, which will be instrumental for computing $Z(n, k; \ell)$ as well (see next section).

Note first that $Z(n, k; 0)$ is simply the partition function of $n + 1$ non-intersecting Dyck paths starting at points $(-2i, 0)$ and ending at points $(0, 2k + 2i)$, $i = 0, 1, ..., n$, as the final portion of the left most path is simply a straight line from $(2k + n, n)$ to $(2n + 2k, 0)$. By applying the Gessel–Viennot formula of lemma 2.2, we have:

**Theorem 4.1.**

$$Z(n, k; 0) = \det_{i,j \in [0,n]} (c_{k+i+j}) = \prod_{i=0}^{n} U_{ij}, \quad U_{ij} = \frac{(2i + 1)!(2k + 2i)!}{(k + 2i + 1)!(k + 2i)!}.$$  \hfill (4.5)

where $c_m = \frac{1}{m+1} \binom{2m}{m}$ is the $m$th Catalan number, enumerating the Dyck paths of $2m$ steps.
Proof. By $LU$ decomposition of the matrix $A$ with entries $A_{ij} = c_{k+i+j}$. Introduce the lower triangular matrix $L$ with entries:

$$L_{ij} = 2^{i-j} \binom{i}{j} \prod_{s=j+1}^{i} \frac{2k+2s-1}{k+1+j+s}$$

$$= \frac{(2k+2i)!k!k+2j+1)!}{(2k+2j)!k!(k+i+j+1)!} \binom{i}{j} = \frac{(2i+2j)!}{(i-j)!} \frac{(2i-j)!}{(i-j)!} \binom{i}{j}.$$ (4.6)

Then $L^{-1}$ has entries:

$$(L^{-1})_{ij} = (-2)^{i-j} \binom{i}{j} \prod_{s=j+1}^{i} \frac{2k+2s-1}{k+i+s} = (-1)^{i+j} \frac{(2k+2i)!k!k+i+j)!}{(2k+2j)!k!(k+i+j+1)!} \binom{i}{j}.$$ (4.7)

Indeed, we have:

$$\sum_{r=j}^{i} (L^{-1})_{ij} L_{ij}$$

$$= \sum_{r=j}^{i} (-1)^{i+r} \frac{(2k+2i)!k!(k+r)!}{(2k+2r)!k!(k+i+r)!} \frac{r}{(2k+2j)!k!(k+i+j+1)!} \binom{i}{j} \frac{(2k+2r)!k!(k+j+1)!}{(2k+2j)!k!(k+i+j+1)!} \binom{j}{r}$$

$$= \frac{(2k+2i)!k!(k+j)!}{(2k+2j)!k!(k+i+j+1)!} \sum_{r=j}^{i} (-1)^{i+r} \frac{(k+i+r)!}{(k+j+r+1)!} \frac{(k+i+j+1)!}{(k+i+j)!} \frac{(k+r)!}{(i-r)!(r-j)!}$$

$$= \frac{1}{(k+i+j+1)} \sum_{r=j}^{i} (-1)^{i+r} \binom{k+i+r}{k+i+j+1} \binom{k+i+j+1}{k+i+1} \binom{k+i+j+1}{k+i} \binom{k+i+j+1}{k+i}$$

If $i = j$, the result is obviously 1, as only the value $r = i = j$ contributes. If $i > j$, the sum reads:

$$(-1)^{i+j} \sum_{r=0}^{i-j} (-1)^{r} \frac{(k+i+j+r)!}{(k+i+1)!} \frac{(k+i+1)!}{(1+x)^{k+i+j+r+1}} \bigg|_{x=-i} = (-1)^{i+j} \frac{(1+x)^{k+i+j+1}}{(1+x)^{k+i+j+1}} \bigg|_{x=-i} = 0$$

by direct application of lemma 2.3.

Finally, let us compute for $i > j$:

$$(L^{-1}A)_{ij} = \sum_{r=0}^{i-j} (-1)^{i+r} \frac{(2k+2i)!k!(k+r)!}{(2k+2r)!k!(k+i+r)!} \frac{r}{(2k+2j)!k!(k+i+j+1)!} \binom{i}{j} \frac{(2k+2r)!k!(k+i+j+1)!}{(2k+2j)!k!(k+i+j+1)!} \binom{j}{r}$$

$$= \frac{(2k+2i)!k!(k+i)!k!(2j)!}{(k+i)!(k+2j)!(k+i+j+1)!} \times \sum_{r=0}^{i-j} (-1)^{i+r} \binom{k+i+r}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1} \binom{k+i+1}{k+i+1}$$.
We now use the following.

**Lemma 4.2.** For all \( i \geq j \), we have:

\[
\begin{align*}
\frac{\partial}{\partial r} f_{k,i,j} := \sum_{r=0}^{i} (-1)^{i-r} \binom{k + r + i}{i-r} \binom{k + 2r + 2j}{2j} &= \binom{(2i+1)(k+i)}{2i+j} & \text{if } i = j, \\
&= 0 & \text{if } i > j.
\end{align*}
\]

**Proof.** Let us write \( j = i - m \). Then:

\[
\begin{align*}
f_{k,i,j-m} &= \sum_{r=0}^{i} (-1)^{i-r} \binom{k + r + i}{i-r} \binom{k + 2i - m + 1}{2i - 2m} \binom{k + 2r + 2j}{2j} \\
&= \frac{1}{m!} \frac{d}{d^m} \left[ \sum_{i=0}^{m} \left( \frac{1 - x^2}{1 - y^2} \right)^{2i+2m} \left( \frac{x}{1 - y^2} \right)^{k+2r+2j} \right]_{r=m=1,
\end{align*}
\]

where in the second line we have used lemma 2.5 to represent \( \binom{k+r+i}{k} \) and \( \binom{k+r}{k} \) and lemma 2.4 to represent the other two binomial coefficients, and where in the third line we performed a change of variable \( x \to x^2 \). Let us now use lemma 2.3 to write:

\[
\begin{align*}
\sum_{i=0}^{m} (tu)^{-i} \left( \frac{1 - x^2}{1 - y^2} \right)^{2i} \left( \frac{x}{1 - y^2} \right)^{k+2r+2j} &= \frac{1}{(1 - x^2)^{2i+2m}} \left( \frac{x}{1 - y^2} \right)^{k+2r+2j} \\
&\text{as the derivatives w.r.t. } u \text{ vanish as soon as } r < 0 \text{ (the overall power of u is } k + r < k). \text{ We may therefore rewrite}
\end{align*}
\]

\[
\begin{align*}
f_{k,i,j-m} &= \frac{1}{m!} \frac{d}{d^m} \left[ \sum_{i=0}^{m} \left( \frac{1 - x^2}{1 - y^2} \right)^{2i} \left( \frac{x}{1 - y^2} \right)^{k+2r+2j} \right]_{r=m=1}.
\end{align*}
\]

It is clear that the quantity:

\[
\begin{align*}
\frac{1}{m!} \frac{d}{d^m} \left[ \sum_{i=0}^{m} \left( \frac{1 - x^2}{1 - y^2} \right)^{2i} \left( \frac{x}{1 - y^2} \right)^{k+2r+2j} \right]_{r=m=1}
\end{align*}
\]
is a polynomial of \( x \), as the denominator after setting \( u = t = 1 \) divides the numerator. The degree of this polynomial is \( 2(k + 2i - m + 1) - (2i - 2m + 1) = 2k + 2i + 1 \), and we must pick the coefficient of \( x^{2k+2i} \) in this polynomial. This coefficient is easily found by expanding equation (4.9) around \( x = \infty \): \[ x^{k+2i}u^{k+1}(1 - \frac{\ell}{m})^{k+2i+1} = (-1)^{k+m}x^{k+2i}(tu)^{k+2i+1} \left( 1 + \frac{2i - 2m + 1}{x} + O(x^{-2}) \right) \]

and we conclude that \[ f_{k,i,m} = \frac{1}{m!} \frac{d^m}{du^m} \left( \frac{1}{k!} \frac{d^k}{du^k} (-1)^k(2i + 1)u^{-i-1} \right)_{u=1} = (2i + 1) \left( \frac{k + i}{k} \right). \]

For \( m = 0 \) this gives \[ f_{k,i,0} = 1! \frac{d^1}{du^1} (-1)^k(2i + 1)u^{-i-1} \big|_{u=1} = (2i + 1) \left( \frac{k + i}{k} \right), \]

while for \( m > 0 \) we find: \( f_{k,i,m} = 0 \) as the derivative w.r.t. \( t \) vanishes. The lemma follows.

We conclude that \( L^{-1}A = :U \) is upper triangular, with diagonal elements:

\[ U_{ij} = \frac{(2k + 2i)!i!(2i)!}{(k + i)!(k + 2i)!} \times \frac{(k + i)!}{i!k!} \]

\[ = \frac{(2i + 1)!}{(k + 2i + 1)} \times \frac{(2k + 2i)!}{(k + 2i)!} \]

and the theorem follows.

\[ \square \]

### 4.3. One-point function

We now turn to the computation of \( H(n,k;\ell) \).

The partition function \( Z(n,k;\ell) \) for general \( \ell \in [0,n+1] \) corresponds by the Gessel–Viennot formula to the determinant of the matrix \( C(n,k,\ell) \) with entries:

\[ C(n,k,\ell)_{ij} = \begin{cases} c_{k+i+j} & \text{if } j \in [0,n-1] \\ b_{2k+2i+n-\ell,n+\ell} & \text{if } j = n \end{cases} \quad (4.10) \]

Indeed, as before \( c_{k+i+j} \) is the number of Dyck paths from \((-2i,0)\) to \((2k+2i,0)\), while \( b_{2k+2i+n-\ell,n+\ell} \) is the number of Dyck paths from \((-2i,0)\) to \((2k+n-\ell,n+\ell)\), where for \( h = a \mod 2 \), we have:

\[ b_{u,h} = \left( a \cdot \frac{a}{a+h} \right) - \left( a \cdot \frac{a-h}{2} - 1 \right) = \frac{h + 1}{a+h+1} \left( a \cdot \frac{a}{a+h} \right) \quad (4.11) \]

so that

\[ b_{2k+2i+n-\ell,n+\ell} = \frac{n + \ell + 1}{n + i + k + 1} \left( \frac{2k + 2i + n - \ell}{k + i - \ell} \right). \quad (4.12) \]

Note also that for \( \ell = 0 \), \( b_{2k+2i+n,0} \neq c_{n+i+k} \), yet the determinant is the same as that with \( \{c_{n+i+k}\}_{i=0}^{n} \) as the last column. This is a simple manifestation of the fact that non-intersecting paths ending at \((2k+2i,0)\), \( j = 0, 1, ..., n \) must end with a number of descending steps,
which is \( j \) for the \( j \)th path counted from the bottom, hence in particular we may cut the topmost path to its last visited vertex before the \( n \) forced descents, namely at the point \((2k + n, n)\).

We have therefore

\[
Z(n, k; \ell) = \det_{ij \in [0, n]} \left( \hat{A}_{ij} \right) = \begin{pmatrix}
  c_k & c_{k+1} & \cdots & c_{n+k-1} & b_{2k+n-\ell, n+\ell} \\
  c_{k+1} & c_{k+2} & \cdots & c_{n+k} & b_{2k+n-\ell+1, n+\ell} \\
  \vdots & \vdots & \ddots & \vdots & \vdots \\
  c_{n+k} & c_{n+k+1} & \cdots & c_{2n+k-1} & b_{2k+3n-\ell, n+\ell}
\end{pmatrix},
\]

(4.13)

To compute this determinant, we use the result of section 2.2. Using the \( L \) matrix of (4.6), we get that \( L^{-1} \hat{A} = \tilde{U} \) where \( \tilde{U} \) is an upper triangular matrix differing from \( U \) only in its last column, with in particular:

\[
\tilde{U}_{n,n} = \sum_{r=0}^{k} (L^{-1})_{n,n} b_{2k+n-\ell+2r, n+\ell} = \sum_{r=0}^{n} (-1)^{n+r} \frac{(2n+2k)!(k+r)!(n+k+r)!}{(2k+2r)!(n+k)!(2n+k)!} \binom{n}{r} \binom{n+\ell+1}{n+r+k+1} \binom{2k+2r+n-\ell}{k+r-\ell}.
\]

As explained in section 2.2, the one-point function \( H(n, k; \ell) \) is nothing but the quantity

\[
H(n, k; \ell) = \frac{\tilde{U}_{n,n}}{U_{n,n}} = \frac{(k+2n+1)!}{(2n+1)!(n+k)!} \times \sum_{r=0}^{n} (-1)^{n+r} \frac{(k+r)!(n+k+r)!(2k+2r+n-\ell)!}{(2k+2r)!(n-r)!(n+r+k+1)!(k-r)!}.
\]

We have the following.

**Theorem 4.3.** The one-point function \( H(n, k; \ell) = \frac{\tilde{U}_{n,n}}{U_{n,n}} \) reads:

\[
H(n, k; \ell) = \frac{(n+\ell+1)!(2k+n-\ell)!}{(2n+2k)!} \sum_{s=0}^{n} \frac{(2n+k-s)!}{(\ell+2s-n)!(n+k-\ell-s)!(2n+1-2s)!} = \frac{1}{(2n+2k)!} \sum_{s=0}^{n} \frac{(n+\ell+1)!(2n+k-s)!}{(2n+1-2s)!} \binom{2n+k-s}{n+\ell}.
\]

(4.14)

In the remainder of this section the theorem is proved in two steps, as we must distinguish the cases \( n \geq \ell \) and \( n < \ell \).

**4.3.1. Case \( n \geq \ell \).** We first note that for \( n \geq \ell \), we may rewrite:

\[
\tilde{U}_{n,n} = \frac{(n+\ell+1)}{(k+2n)!(n-k)!} \frac{(2n+2k)!}{(n+\ell)!(n-k)!} \times \sum_{r=0}^{n} (-1)^{n+r} \binom{n+k+r}{k+r} \binom{2k+2r+n-\ell}{k-\ell} \binom{2k+2n+1}{n-r}.
\]

We have the following:
Lemma 4.4. For all $k \geq 0$ and all $n \geq \ell \geq 0$, we have:

$$g_{n,k,\ell} := \sum_{r=0}^{n} (-1)^{n+r} \binom{n+k+r}{n+k} \frac{(k+r)}{l} \frac{(2k+2r+n-\ell)}{n-\ell} \frac{(k+2n+1)}{n-r} = \left(\frac{2n+1}{n-\ell}\right) \binom{n+\ell}{n}.$$  

(4.15)

Proof. We compute:

$$g_{n,k,\ell} = \frac{1}{(n+k)!} \frac{d^{n+k} 1}{d t^{n+k}} \frac{d^{\ell} \ell!}{d u^{\ell}} \sum_{r=0}^{n} t^{\ell + k + r} u^{k+r} \left(\frac{(1-x)^{k+2n+1}}{(1-x)^{n-\ell+1}}\right)\bigg|_{t=u=1}.$$  

As before, we wish to use lemma 2.3, which in this case gives:

$$\sum_{r=-k}^{n} (tu)^{\ell-n} \frac{(1-x^{2})^{k+2n+1}}{(1-x)^{n-\ell+1}}\bigg|_{x^{2n-2\ell+2+n}} = \frac{(1-x^{2})^{k+2n+1}}{(1-x)^{n-\ell+1}}\bigg|_{x^{2n+2}}.$$  

As before, the discrepancy between the summation ranges is resolved by noting that terms with $r < 0$ contribute 0 after taking $n + k$ derivatives w.r.t. $t$, as the corresponding overall power of $t$ is $n + k + r < n + k$. We therefore have:

$$g_{n,k,\ell} = \frac{1}{(n+k)!} \frac{d^{n+k} 1}{d t^{n+k}} \frac{d^{\ell} \ell!}{d u^{\ell}} \sum_{r=0}^{n} t^{\ell-k} u^{k} \frac{\ell-k}{\ell} \frac{(1-tu^{2)})^{k+2n+1}}{(1-tu)^{n-\ell+1}}\bigg|_{t=u=1}.$$  

As before it is readily seen that the quantity

$$\frac{1}{(n+k)!} \frac{d^{n+k} 1}{d t^{n+k}} \frac{d^{\ell} \ell!}{d u^{\ell}} t^{-k} u^{-n-k} \frac{(1-tu^{2})^{k+2n+1}}{(1-tu)^{n-\ell+1}}\bigg|_{t=u=1}$$

is a polynomial of $x$ as the denominator at $t = u = 1$ always divides the numerator. The degree of this polynomial is $2(k+2n+1) - (n-\ell+1) = 2k+3n+\ell+1$. The desired result is obtained by an expansion around $x \to \infty$:

$$t^{-k} u^{-n-k} \frac{(1-tu^{2})^{k+2n+1}}{(1-tu)^{n-\ell+1}} = (-1)^{k+\ell} \sum_{i=0}^{k+2n+1} \sum_{j=0}^{\ell+i-j} t^{\ell-i-j} u^{\ell-i-j} \frac{k+2n+1}{i} \frac{n+\ell+j}{n-\ell} x^{2k+3n+\ell+1+2n+1-\ell-j}.$$
and extracting the coefficient of $x^{2n+2k}$, namely imposing $2i + j = n + \ell + 1$. This yields:

$$g_{n,k,\ell} = (-1)^{n+k+\ell} \frac{1}{(n+k)!} \frac{1}{d^{n+k+\ell}!} \frac{d^n}{du^n} \sum_{i=0}^n i^{-1} u^{i-n-1} \binom{k + 2n + 1}{i} \binom{2n + 1 - 2i}{n - \ell} \bigg|_{r=u=1}.$$ 

Note that each term with $i \geq 1$ yields zero after differentiation with respect to $\ell$, as $i - 1 < n < k + n$. We are left with the contribution of $i = 0$, which reads:

$$g_{n,k,\ell} = (-1)^{n+k+\ell} \frac{2n + 1}{(n+k)!} \frac{1}{d^{n+k+\ell}!} \frac{d^n}{du^n} \bigg|_{r=u=1} = \binom{2n + 1}{n - \ell} \binom{n + \ell}{\ell}$$

and the lemma follows.

We deduce that for $n \geq \ell$, we have:

$$\frac{\tilde{U}_{n,n}}{U_{n,n}} = \frac{(n + \ell + 1) (2n + 2k)! (n! \ell!)^2 (2n + 1) \binom{n + \ell}{\ell}}{(2n + 1)! (2n + 2k)!} = 1.$$ 

To identify this with the statement of theorem 4.3 for $n \geq \ell$, we need to show that equation (4.14) gives $H(n; k; \ell) = 1$ for $n \geq \ell \geq 0$. This is a consequence of the following:

**Lemma 4.5.** The following identity holds for all $n \geq \ell \geq 0$:

$$\sum_{s=0}^n \binom{n + \ell + 1}{2n + 1 - 2s} \binom{2n + k - s}{n + \ell} = \binom{2n + 2k}{n + \ell}.$$ 

**Proof.** We write:

$$\sum_{s=0}^n \binom{n + \ell + 1}{2n + 1 - 2s} \binom{2n + k - s}{n + \ell} = \sum_{s=0}^n \binom{1 + x}^{n + \ell + 1} \bigg|_{x^s} \frac{1}{1 - y^{n + \ell + 1}} \bigg|_{y^{n + \ell + 1}}$$

$$= \sum_{s=0}^n \binom{1 + x}^{n + \ell + 1} \bigg|_{x^s} \frac{1}{1 - y^{n + \ell + 1}} \bigg|_{y^{2(n + \ell + 1)}}$$

$$= \left. \left( \frac{1 + x}{1 - x^2} \right)^{n + \ell + 1} \right|_{x^{2n + \ell}} = \binom{2n + 2k}{n + \ell}$$

where we have performed a change of variables $y \rightarrow y^2$ in the second line, and where to go from the second to the third line we noticed that the summation over $s$ can be relaxed to all $s \in [0, n + \ell + 1]$, as the actual range of the contributing terms is for $s \in \left[ \frac{n + \ell}{2}, n \right] \subset [0, n + \ell + 1]$, and applied lemma 2.3 (note that this would be wrong if $\ell > n$). The result of the lemma follows.

This completes the proof of theorem 4.3 in the case $n \geq \ell$. 


4.3.2. Case $\ell > n$. We now consider the case $\ell > n$.

We may write:

$$\frac{\bar{U}_{n,n}}{U_{n,n}} = \frac{(n + \ell + 1)^{n\ell!}}{(2n + 1)!} \times \sum_{r=0}^{n} (-1)^{n+r} \binom{n + k + r}{n+k} \binom{k + r}{l} \binom{k + 2n + 1}{n-r} \frac{(2k + 2r + n - \ell)!}{(2k + 2r)!}.$$  \hspace{1cm} (4.16)

In the following, we use the standard notation for Pochhammer symbols:

$$(x)_m = x(x - 1) \cdots (x - m + 1).$$

Introduce the following quantities, using respectively the expression of $\frac{\bar{U}_{n,n}}{U_{n,n}}$ via equation (4.16) and that of the sought after result for $H(n,k;\ell)$ via equation (4.14):

$$P_{n,\ell}(k) := \frac{(2n + 1)![(n + \ell)!\binom{2n + 2k}{n + \ell}]}{(n + \ell + 1)!\binom{n + \ell}{n + \ell}^2} \frac{\bar{U}_{n,n}}{U_{n,n}}$$

$$= \sum_{r=0}^{n} (-1)^{n+r} \binom{n + k + r}{n+k} \binom{k + r}{l} \binom{k + 2n + 1}{n-r} \frac{(2k + 2r + n - \ell)!}{(2k + 2r)!} \cdot$$

and

$$Q_{n,\ell}(k) := \frac{(2n + 1)![(n + \ell)!\binom{2n + 2k}{n + \ell}]}{(n + \ell + 1)!\binom{n + \ell}{n + \ell}^2} H(n,k,\ell)$$

$$= \frac{(2n + 1)!}{n + \ell + 1} \sum_{r=0}^{n} \binom{n + \ell + 1}{2n + 1 - 2s} (k + 2n - s)_{n-s} (n + k - \ell)_s.$$  \hspace{1cm} (4.18)

$P_{n,\ell}$ and $Q_{n,\ell}$ are two polynomials of $k$, with a priori degree $\deg(P_{n,\ell}) = 3n$ and $\deg(Q_{n,\ell}) = n$, but we have the following:

**Theorem 4.6.** The polynomials $P_{n,\ell}$ and $Q_{n,\ell}$ have same degree $n$. Moreover, they coincide on all the points $k = -j - n$, for $j = 1, 2, ..., n + 1$, where they take the values:

$$P_{n,\ell}(-j - n) = Q_{n,\ell}(-j - n) = (-1)^n \binom{2n + 1}{n + \ell + 1}! [(j + 1)! (n + 2j + \ell + 1)! (n + \ell + 1)! (2j - 1)! (\ell + j - 1)!].$$  \hspace{1cm} (4.19)

**Proof.** Let us first compute the degree of $P_{n,\ell}$. First note that the expression (4.17) is valid for both $\ell > n$ and $\ell \leq n$. But in the latter case, we have found that

$$\frac{\bar{U}_{n,n}}{U_{n,n}} = \frac{(n + \ell + 1)!\binom{n + \ell}{n + \ell}^2}{(2n + 1)!\binom{n + \ell}{n + \ell}^2} P_{n,\ell}(k) = 1$$

which implies that
To conclude, we need a variant of lemma 4.5, proved in the same way:

is a polynomial of degree \( n \), as for \( n \geq \ell \) the denominator always divides the numerator. We may infer that all the cancellations that reduce the degree from \( 3n \) to \( n \) still occur when \( \ell > n \) and the degree of \( P_{n,\ell} \) is always \( n \). Let us now compute using (4.17) for \( k = -j - n, 1 \leq j \leq n + 1 \):

\[
P_{n,\ell}(-j - n) = \sum_{i=0}^{n} (-4)^{n-i} \binom{n}{i} (2r - 2j - n - \ell)_{2r} (r - j)_{r} (n - j + 1)_{n-j} (-j - \frac{1}{2})_{n-r}.
\]

Vanishing contributions to the sum arise from the terms \((r - j)\), if \( r \geq j \) and \((n - j + 1)_{n-r}\) if \( r < j - 1 \). We conclude that only the term \( r = j - 1 \) survives, with value:

\[
P_{n,\ell}(-j - n) = (-4)^{n-j+1} \binom{n}{j-1} (2r - 2j - n - \ell)_{2j-2}
\]

\[
\times (-1)_{j-1} (-j - \ell)_{n-j+1} (n - j + 1)_{n-j+1} (-j - \frac{1}{2})_{n-j+1}
\]

\[
= (-2)^{n-j+1} \binom{n}{j-1} (2j + n - \ell - 1)! \frac{(n + \ell)!}{(n + \ell + 1)!}
\]

\[
\times (-1)^{j-1} (j - 1)! (n - j + 1)! (n + \ell)! (2n + 1)! (2j - 1)!
\]

\[
= (-1)^{n} (2n + 1)! (j - 1)! (2j + n - \ell - 1)! (n + \ell + 1)! (j + \ell - 1)!
\]

which proves (4.19) for \( P_{n,\ell} \). Finally, let us compute using (4.18) for \( k = -j - n, 1 \leq j \leq n + 1 \):

\[
Q_{n,\ell}(-j - n) = \frac{(2n + 1)!}{n + \ell + 1} \sum_{s=0}^{n} \binom{n + \ell + 1}{2n + 1 - 2s} (n - j - s)_{n-s} (-j - \ell)_{s}.
\]

The vanishing contributions to the sum correspond to \( n - j \geq s \), hence we may rewrite:

\[
Q_{n,\ell}(-j - n) = (-1)^{n} \frac{(2n + 1)!}{n + \ell + 1} \sum_{s=n-j+1}^{n} \binom{n + \ell + 1}{2n + 1 - 2s} (j - 1)! (j + \ell + s - 1)!
\]

\[
\times (j + \ell - 1)!
\]

\[
= (-1)^{n} \frac{(2n + 1)!}{n + \ell + 1} \sum_{s=n-j+1}^{n} \binom{n + \ell + 1}{2n + 1 - 2s} (j + \ell + s)! (j + \ell + s - 1)!
\]

\[
\times (j + \ell - 1)!
\]

\[
= (-1)^{n} \frac{(2n + 1)!}{n + \ell + 1} \sum_{s=n-j+1}^{n} \binom{n + \ell + 1}{2n + 1 - 2s} (j + \ell + s - 1)!
\]

\[
\times \frac{(j + \ell + s - 1)!}{(j + \ell + s - 1)!}
\]

\[
\times \frac{(2n + 1)!}{2n + 1 - 2s} (j + \ell - 1)!
\]

To conclude, we need a variant of lemma 4.5, proved in the same way:
Lemma 4.7. We have the identity:

\[
\frac{1}{(2^{s+n+\ell-1})} \sum_{j=n-j+1}^{n} \binom{n+\ell+1}{n+\ell} \binom{j+\ell+s-1}{n+\ell} = 1
\]

for \( \ell > n \) and \( 1 \leq j \leq n+1 \).

This is proved by computing:

\[
\sum_{j=n-j+1}^{n} \binom{n+\ell+1}{n+\ell} \binom{j+\ell+s-1}{n+\ell} = \sum_{j=n-j+1}^{n} \frac{(1+x)^{n+\ell+1}}{(1-x^2)^{n+\ell+1}} \left| \frac{1}{2^{j+s-1}} \right| \left| \frac{j+\ell+s}{2} \right|
\]

by application of lemma 2.3. We deduce that (4.19) holds for \( Q_{n,\ell} \) as well, and theorem 4.6 follows.

This completes the proof of theorem 4.3 in the case \( \ell > n \).

4.4. Tangent method and arctic curve

Using the results of previous section, we are now ready to apply the saddle-point approximation to the normalized sum:

\[
\frac{N(n, k, p)}{Z(n, k; 0)} = \sum_{\ell=0}^{n+k} H(n, k; \ell) Y(p; \ell).
\]

The tangent method described in section 2.1, applied to the present problem, is illustrated in figure 4. It consists of finding which value of \( \ell \) dominates the contribution to the sum (4.20), by the saddle-point method. The corresponding parametric family of lines \( CD \) is tangent to the arctic curve, which is uniquely defined as their envelope.

We wish to evaluate the quantity \( N(n, k, p)/Z(n, k; 0) \) at large \( n, k, p \) via the saddle-point method. Write \( k = xn, p = yn, n \) large. The summation in (4.20) is approximated by the integral

\[
\frac{N(n, xn, yn)}{Z(n, xn; 0)} \approx n \int_{0}^{1+\xi} d\xi H(n, xn; \xi n) Y(yn; \xi n)
\]

by writing \( \ell = \xi n \) and \( \sum_{\ell=0}^{n+k} \rightarrow n \int_{0}^{1+\xi} d\xi \). For large \( n \), using the expression (4.4), we have by the Stirling formula:

\[
Y(yn; \xi n) \approx \left( \frac{n(y + \xi)}{n\xi} \right) e^{S_0(y, \xi)}
\]

where

\[
S_0(y, \xi) = (y + \xi) \log(y + \xi) - \xi \log(\xi) - y \log(y).
\]

Similarly, we write the summation (4.14) as:
组装所有的部分，我们得到鞍点方程

\[ S(x, y, \xi, \sigma) = S_0(y, \xi) + S_1(x, \xi, \sigma), \]

其中变量 \( \xi, \sigma \) 必须满足

\[ 0 \leq \sigma \leq 1, \quad 0 \leq \xi \leq 1 + x, \quad \xi + \sigma \leq 1 + x, \quad \xi + 2\sigma \geq 1. \]  \( \text{(4.24)} \)

我们找到:

\[ \partial_{\xi} S = 0 = \log \left( \frac{(y + \xi)(1 + \xi)(1 + x - \xi - \sigma)}{\xi(2x + 1 - \xi)(\xi + 2\sigma - 1)} \right) \]

\[ \partial_{\sigma} S = 0 = \log \left( \frac{(1 - \sigma)^2(1 + x - \xi - \sigma)}{(\xi + 2\sigma - 1)^2(2 + x - \sigma)} \right). \]
The second equation boils down to:
\[
\sigma = \frac{(2 + 3x) - \xi(2 + x)}{3 + 4x - \xi}.
\]

We see that all conditions (4.24) are automatically satisfied (as \(x, y > 0\)) except for the bound \(\sigma \geq 0\). Define
\[
\xi^* := \frac{2 + 3x}{2 + x} = 1 + \frac{2}{2 + x} x \in (1, x + 1).
\]

If \(\xi < \xi^*\): then \(\sigma > 0\), and the first equation gives:
\[
y(2x + 1 - \xi^2)(1 + \xi) = 0
\]
which has no physical solution. Hence we must have: \(\xi > \xi^*\), but then the second equation cannot hold, which means that the variable \(\sigma\) is saturated to be \(\sigma = 0\). In that case, we end up with the first equation at \(\sigma = 0\), namely a cubic relation for \(\xi\):
\[
A(x, y, \xi) := (y + \xi)(1 + \xi)(1 + x - \xi) - \xi(2x + 1 - \xi)(\xi - 1) = 0
\]
with a unique solution \(\xi(x, y)\) in the physical domain (4.24).

We note that \(A(x, 0, \xi) = \xi((2 + 3x) - (2 + x)\xi)\), with the only admissible solution \(\xi = \xi^*\). We deduce that the arctic curve is tangent to the right boundary line (with equation \(v = 2x + u\)) in the physical domain (4.24).

Similarly, for large \(y\) we have \(\lim_{y \to \infty} y^{-1}A(x, y, \xi) = (1 + \xi)(1 + x - \xi)\), with only admissible solution \(\xi = 1 + x\). We deduce that the arctic curve is tangent to the left boundary (with equation \(v = 2x + u\)) at some point \(A\) and \(B\) are represented in figure 4 in scale of \(n\).

Assuming the tangent to the sought-after arctic curve is the parametric line \((CD)\): 
\(v = a(x, y)u + b(x, y)\), with \(\xi = \xi(x, y)\), we have
\[
1 + y = (2x + 1 + y)a + b, \quad 1 + \xi = (2x + 1 - \xi)a + b
\]
hence
\[
a = \frac{y - \xi(x, y)}{y + \xi(x, y)}, \quad b = \frac{2(1 + x + y)\xi(x, y) - y}{y + \xi(x, y)}.
\]

Note that \(A(x, y, \xi) = 0\) may be inverted as
\[
y = y(x, \xi) := \frac{\xi((2 + x)\xi - (2 + 3x))}{(1 + \xi)(1 + x - \xi)}.
\]

The envelope of the parametric family of lines \(F(u, v, \xi, x) := v - a(x, y(\xi))u - b(x, y(\xi))\) is the solution of the system
\[
F(u, v, \xi, x) = 0 \quad \text{and} \quad \partial_\xi F(u, v, \xi, x) = 0.
\]
Eliminating \(\xi\) between these equations, we find the ellipse:
\[ x^2 v^2 - 4(1 + x)u(2x - u) = 0. \]  \hspace{1cm} (4.25)

Strictly speaking, we have only computed the portion of arctic curve between the two tangency points \( A \) and \( B \), namely for \( u \in \left[ \frac{2x}{1+x}, \frac{2(1+\frac{1}{x})}{2+\frac{1}{x}} \right] \). The entire half-ellipse with \( v \geq 0 \) turns out to be the full arctic curve for our problem, as shown in next section.

5. An equivalent rhombus tiling problem

In this section we present an equivalent rhombus tiling problem to that of non-intersecting Dyck paths of the previous section, and use it to derive the remaining portions of the arctic ellipse (4.25).

5.1. The equivalent tiling problem

The problem studied so far has an interpretation in terms of rhombus tilings of a ‘half-hexagon’, as illustrated in figure 5. Such a tiling is entirely determined by either of three families of non-intersecting paths, which connect two parallel edges in each tile (the three families of paths are represented in black, blue and red respectively).

We are now in position to complete the previous study of the arctic curve for the Dyck path problem. We simply have to apply the tangent method for the other two families of paths. We concentrate on the red paths, as the treatment of the blue paths is completely identical, up to reflection w.r.t. to the vertical line through \((k, 0)\).

5.2. Exact enumeration

We may now consider the red path family depicted in figure 6 that has starting points at \((i, i + 2n + 2), i = 0, 1, ..., k - 1\) and endpoints at \((2j, 0), j = 0, 1, ..., k - 2\) as well as the remote point \((2k - 2 + p, -p)\). Note that the vertical steps are \((0, -2)\) and the diagonal ones are \((1, -1)\). Let us denote by \(N_2(n, k, p)\) the total number of path configurations in this family. As before, we introduce the numbers \(Z_2(n, k; \ell)\) and \(Y_2(p; \ell)\) that count respectively the non-intersecting family from \((i, i + 2n + 2), i = 0, 1, ..., k - 1\) to \((2j, 0), j = 0, 1, ..., k - 2\) as well as \((2k - 2, 2\ell), \ell \geq 0\), and the configurations of a single path from \((2k - 1, 2\ell - 1)\) to \((2k - 2 + p, -p)\). We have the decomposition formula:

\[ N_2(n, k, p) = \sum_{\ell = 0}^{n+1} Z_2(n, k; \ell) Y_2(p; \ell). \]  \hspace{1cm} (5.1)

The latter number is easily found to be:

\[ Y_2(p; \ell) = \binom{p + \ell - 1}{\ell} = Y(p; \ell) \]

as the path has \( p - 1 \) diagonal and \( \ell \) vertical steps.

The number \(Z_2(n, k; 0)\) is given via Gessel–Viennot by the following determinant:
Theorem 5.1. The number $Z_2(n, k;0)$ is given by:

$$Z_2(n, k;0) = \det(A), \quad A_{ij} = \binom{j+n+1}{2j-i}, \quad i, j = 0, 1, \ldots, k-1.$$ 

Proof. We proceed as before by LU decomposition of $A$. The LU decomposition of $A$ uses the lower triangular matrix $L$ with elements:

$$L_{ij} = \frac{(2i-2j)(i+j)}{(i+2n+2)(i-j)} = \frac{i!(j+2n+2)!}{(i-j)!(2j-i)!(i+2n+2)!},$$

and the entries of its inverse read:

$$(L^{-1})_{ij} = (-1)^{i+j} \frac{(j+2n+2)!(i-1)!(2i-j-1)!}{(i+2n+2)!(j-1)!} \binom{j+n+1}{2j-i} = (-1)^{i+j} \frac{(j+2n+2)!(2i-j-1)!}{(i+2n+2)!(i-j)!(j-1)!}.$$ 

This is readily checked by computing for $i \geq j$:

$$\sum_{m=0}^{k-1} (L^{-1})_{im}L_{mj} = \sum_{m=0}^{k-1} (-1)^{i+m} \frac{(m+2n+2)!(2i-m-1)!}{(i+2n+2)!(i-m)!(m-1)!(m-j)!(2j-m)!(m+2n+2)!} = (-1)^{i+j} \frac{(j+2n+2)!}{(i+2n+2)!} \sum_{m=i}^{\text{Min}(j,n)} m(-1)^m \frac{(2i-m-1)!}{(2j-m)!(i-m)!(m-j)!}.$$ 

If $i = j$, this is clearly equal to 1 as only $m = i = j$ contributes. If $i > j$, we have:
\[
\frac{(i-j)!}{(2i-2j-1)!} \sum_{m=j}^{\text{Min}(i,2j)} m(-1)^m \frac{(2i-m-1)!}{(2j-m)!}(i-m)! (m-j)!
\]

\[
= \sum_{m=j}^{\text{Min}(i,2j)} m(-1)^m \binom{2i-m-1}{2i-2j-1} \binom{i-j}{m-j}
\]

\[
= (-1)^j \sum_{m=j}^{\text{Min}(i,2j)} x^{-2j}(1-x)^{-(2i-2j)} y^{i(1-y)^{i-j}}
\]

\[
= (-1)^j x^{-2j}(1-x)^{-(2i-2j)} y \frac{d}{dy} (y^{i(1-y)^{i-j}}) \bigg|_{y=\rho}
\]

\[
= (-1)^j y^{2j} \frac{1}{(1-y)^{2i-2j}} y \frac{d}{dy} (y^{i(1-y)^{i-j}}) \bigg|_{y=\rho} = 0
\]

as the series of \( y \) has \( y^{2i-j} \) in factor, and hence has no constant term. Note that we have used lemma 2.3 to go from the fourth to the fifth line above.

We have \( L^{-1}A = U \), where \( U \) is upper triangular. To see this we compute for \( i > j \):
\[ \sum_{m=0}^{k-1} (L^{−1})_{j,0} A_{m,j} = \sum_{m=0}^{k-1} (-1)^{j+m} \frac{(m + 2n + 2)!((2i - m - 1))!}{(i + 2n + 2)!((2j - m)!)(m + n + 1)!} \left( j + n + 1 \right) \]

\[ = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \sum_{m=0}^{k-1} (-1)^{j+m} \frac{(2j - 1)!}{(m - 1)!} \frac{(2i - m - 1)!}{(m + n + 1)!} \left( j + n + 1 \right) \]

\[ = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \sum_{m=0}^{k-1} (-1)^{j+m} \frac{(1 + x)^{2j-1}}{x^y} \frac{1}{y^y (1 + y)^y} \frac{1}{z^z (1 - z)^z} \left( j + n + 1 \right) \]

\[ = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \int \frac{1}{(2\pi)^2} \frac{dydz}{yx} \frac{1}{1 - \frac{2}{\pi}} \frac{1}{x^y} \frac{1}{y^y} \frac{1}{z^z (1 - z)^z} \]

where we noted that the contribution of the term \((xy/z)^2\) in the fourth line vanishes as the corresponding \(y\) residue vanishes (as \(k - i > 0\)). Expressing the residue at \(x = 0\) in terms of that at \(x = z/y\), we get:

\[ \sum_{m=0}^{k-1} (L^{−1})_{j,0} A_{m,j} = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \int \frac{1}{(2\pi)^2} \frac{dydz}{yx} \frac{1}{1 - \frac{2}{\pi}} \frac{1}{x^y} \frac{1}{y^y} \frac{1}{z^z (1 - z)^z} \]

Let us write the above as \(R_{y \rightarrow 0}\) to indicate that it picks up the residue at \(y = 0\). First note that by expressing the residue at \(z = 0\) in terms of that at \(z = 1\), we get the identity:

\[ R_{y \rightarrow 0} = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \int \frac{dydz}{(2\pi)^2} \frac{(y + 1 - z)^{2i-1}}{y^y (1 + y)^y z^z (1 - z)^z} \]

\[ (5.2) \]

Next, let us write: \(R_{y \rightarrow 0} = -R_{y \rightarrow 1} - R_{y \rightarrow \infty}\) by use of the Cauchy theorem. To compute \(R_{y \rightarrow 1}\) we perform the change of variables \(y = -1 - t\) in the original integral, leading to:

\[ R_{y \rightarrow 1} = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \int \frac{dt dz}{(2\pi)^2} \frac{(-t - 1 + z)^{2j-1}}{t^t (1 + t)^t z^z (1 - z)^z} \]

\[ = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \int \frac{dt dz}{(2\pi)^2} \frac{(t + 1 - z)^{2j-1}}{t^t (1 + t)^t z^z (1 - z)^z} = R_{y \rightarrow 0} \]

by comparing with the second expression \((5.2)\) above. We deduce that \(R_{y \rightarrow 0} = -\frac{1}{2} R_{y \rightarrow \infty}\) with

\[ R_{y \rightarrow \infty} = \left( \frac{(n + 1 + j)!^2 (i - 1)!}{(i + 2n + 2)!((2j - 1)!)} \right) \int \frac{1}{(2\pi)^2} \frac{du dz}{u^u (1 + u)^{(1 + u)} z^z (1 - z)^z} \]

where we have performed the change of variables \(y = 1/u\) to express \(R_{y \rightarrow \infty}\) as \(R_{u \rightarrow 0}\). If \(i > j\), the residue at \(u = 0\) vanishes, henceforth \((L^{-1}A)_{ij} = -\frac{1}{2} R_{y \rightarrow \infty} = U_{ij} = 0\), and therefore \(U\) is upper triangular. Finally, if \(i = j\), we find:
\[ U_{ij} = \frac{1}{2} \frac{(n+1+i)!}{(i+2n+2)!} (i-1)! \int \frac{dz}{2\pi i} \left( \frac{1}{z^{i+n+2}} \right) \frac{1}{(1-z)^{i+n+2}} \]
\[ = \frac{((n+1+i)!)(2i+2n+2)!}{(i+2n+2)!((2i)!)^2} \frac{(2n+2+i)!}{(2n+2+i)!((2i)!)^2} \]

and the theorem follows.

\[ \square \]

**Remark 5.2.** Comparing theorems 4.1 and 5.1, we obtain the following identity:
\[ \prod_{i=0}^{n} \frac{(2i+1)!}{(k+2i+1)!} = \prod_{i=0}^{k-1} \frac{(2n+2+i)!}{(2n+2+i)!((2i)!)^2} \]

Note in particular that we have:
\[ U_{k-1,i} = \frac{(2k+2n)!}{(2k+n+1)!(2k-2)!} \]

The number \( Z_2(n, k; \ell) \) is given via the Gessel–Viennot formula by the following determinant:
\[ Z_2(n, k; \ell) = \det(\tilde{A}) = \begin{cases} \binom{j+n+1}{2j-i} & \text{if } j = 0, 1, \ldots, k-2 \\ \binom{k+n-k-\ell}{2k-2-i} & \text{if } j = k-1 \end{cases} \]

As before, let us introduce the one-point function: \( H_2(n, k; \ell) := Z_2(n, k; \ell) / Z_2(n, k; 0) \). Defining the upper triangular matrix \( \tilde{U} \) such that \( \tilde{A} = L\tilde{U} \), we may rewrite following section 2.2:
\[ H_2(n, k; \ell) = \frac{\tilde{U}_{k-1,k-1}}{U_{k-1,k-1}}. \]

We have the following.

**Theorem 5.3.** The one-point function \( H_2(n, k; \ell) \) reads:
\[ H_2(n, k; \ell) = \frac{2}{(2n+2)!} \sum_{s=\ell}^{n+1} \binom{k+n-s}{k-2} \binom{k+n+s-1}{k-2}. \]  

**Proof.** We have:
\[ \tilde{U}_{k-1,k-1} = \sum_{i=0}^{k-1} (L^{-1})_{i,k-1} \tilde{A}_{i,k-1} = \sum_{i=0}^{k-1} (-1)^{i+k-1} \frac{(i+2n+2)!}{(k+2n+1)!} \frac{(2k-i-3)!}{(2k-2-i)!} \binom{n+k-\ell}{2k-2-i} \]
\[ = \frac{(2n+3)!}{(k+2n+1)!} \sum_{i=0}^{k-1} (-1)^{i+k-1} \binom{2n+2+i}{i-1} \binom{2k-i-3}{k-1-i} \frac{n+k-\ell}{2k-2-i} \]

which implies:
\[
\frac{\hat{U}_{k-1,k-1}}{U_{k-1,k-1}} = \frac{2}{(2n+3)} \sum_{i=0}^{k-1} (-1)^i \binom{2n + 2 + i}{i} \binom{2k - i}{k - i} \binom{n + k - \ell}{2k - 2i}
\]

\[
= \frac{2}{(2n+3)} \sum_{i=0}^{n+k-\ell} (-1)^i \binom{2n + 2k - i}{2n + 3} \binom{i - 1}{k - 2} \binom{n + k - \ell}{i}
\]

\[
= \frac{2}{(2n+3)} \int_0^1 \int_0^1 \frac{1}{(2\pi)^2 x^{2n+2} \sqrt{1-x}^3} dx \, dy \, (x + y)^{\ell+k-\ell}. \]

Let us rewrite the latter integral by expressing the residue at \( x = 0 \) in terms of that at \( x = 1 \), by the change of variables \( x = 1 - u \):

\[
\frac{\hat{U}_{k-1,k-1}}{U_{k-1,k-1}} = \frac{2}{(2n+3)} \int_0^1 \int_0^1 \frac{1}{(2\pi)^2 u^{n+4} \sqrt{1-u}^3} \frac{1}{(1-u + y)^{\ell+k-\ell}} \frac{1}{(1-u + y)^{\ell+k-\ell}} \frac{1}{(1-u + y)^{\ell+k-\ell}} \frac{1}{(1-u + y)^{\ell+k-\ell}}
\]

\[
\frac{2}{(2n+3)} \sum_{i=0}^{n+k-\ell} (-1)^i \binom{2n + 2k - i}{2n + 3} \binom{i - 1}{k - 2} \binom{n + k - \ell}{i}
\]

\[
= \frac{2}{(2n+3)} \sum_{i=0}^{n+k-\ell} \binom{k + n - \ell - i - 1}{k - 2} \binom{n + k + \ell + i}{k - 2}
\]

\[
= \frac{2}{(2n+3)} \sum_{i=0}^{n+k-\ell} \binom{k + i - 2}{k - 2} \binom{k + 2n + 1 - i}{k - 2}
\]

which completes the proof of (5.4) up to the change of summation variable \( s = n + 1 - i \), and the theorem follows.

5.3. Tangent method

Taking \( n \to \infty \) with \( k = xn, \ell = \xi n, s = \sigma n \) in (5.4), we find:

\[
H_2(n, xn; \xi n) \simeq n \int_0^1 \frac{1}{\xi} d\sigma \, e^{S'_1(x, \sigma)}
\]

where

\[
S'_1(x, \sigma) = 2 \text{Log}(2) + 2x \text{Log}(2x) - (1 + \sigma) \text{Log}(1 + \sigma) - (1 - \sigma) \text{Log}(1 - \sigma) - 2x \text{Log}(x)
\]

\[
- (2 + 2x) \text{Log}(2 + 2x) + (1 + x + \sigma) \text{Log}(1 + x + \sigma) + (1 + x - \sigma) \text{Log}(1 + x - \sigma).
\]
Similarly, we may use the asymptotic expression for \( Y_2(p; \ell) = Y(p; \ell) \) for \( p = \infty \) of equations (4.21) and (4.22). The integration variables must belong to the domain: \( 0 \leq \xi \leq x \). We now have to perform the saddle-point method for the total action \( S(x, y, \xi, \sigma) = S_0(y, \xi) + S'_1(x, \sigma) \) and the integral:

\[
\int_0^x \int_0^1 d\xi \int_0^1 d\sigma \, e^{\kappa S(x, y, \xi, \sigma)}.
\]

We first write:

\[
\partial_\xi S = 0 = \text{Log} \left( \frac{(1 - \sigma)(1 + x + \sigma)}{(1 + \sigma)(1 + x - \sigma)} \right)
\]

with the unique solution \( \sigma = 0 \). This solution is never selected, and we have the saturation by the lower bound \( \sigma = \xi \). This results in the action \( S'(y, x, \xi) = S_0(y, \xi) + S'_1(x, \xi) \), with the bound \( 0 \leq \xi \leq 1 \). We then write the saddle-point condition:

\[
\partial_\xi S' = 0 = \text{Log} \left( \frac{(y + \xi)(1 - \xi)(1 + x + \xi)}{\xi(1 + \xi)(1 + x - \xi)} \right).
\]

We find that for \( 0 \leq \xi \leq x \), we have

\[
y = y(\xi) := \frac{2x\xi^2}{(1 - \xi)(1 + x + \xi)} \in (0, \infty).
\]

Assuming the tangent to the sought-after arctic curve is the parametric line through \((2x + y, -y)\) and \((2x, 2\xi)\), with equation in the \((u, v)\) plane: \( v = a(x, y)u + b(x, y) \), with \( y = y(\xi) \), and

\[-y = (2x + y)a + b, \quad 2\xi = 2xa + b\]

hence

\[
a(\xi) = -\frac{y(\xi) + 2\xi}{y(\xi)}, \quad b(\xi) = 2\frac{(\xi y(\xi) + 2\xi x + xy(\xi))}{y(\xi)}.
\]

As before, the envelope of the family of lines \( F(u, v, \xi) = v - a(\xi)u - b(\xi) \) is the solution to the system \( F = 0 \) and \( \partial_\xi F = 0 \). We find that \( \xi = \frac{y}{(2x + y)} \) and finally, as \( \xi \leq 1 \):

\[
x^2v^2 - 4(1 + x)u(2x - u) = 0, \quad u \in \left[ \frac{2x(1 + x)}{2 + x}, \frac{2x(1 + x)}{2 + x} \right]
\]

which is nothing but the portion of the arctic ellipse (4.25) corresponding to \( u \in \left[ \frac{2x(1 + x)}{2 + x}, 2 \right] \).

Reflection of the problem w.r.t. the vertical line \( u = 1 \) shows that the remaining portion of the arctic curve is another portion of the same ellipse, corresponding to \( u \in \left[ 0, \frac{2x}{2 + x} \right] \).

6. Another path model

We consider non-intersecting square-lattice paths in the first quadrant, with left and up steps \((-1, 0), (0, 1)\), that start at the points \((2i, 0), i = 0, 1, ..., n\) and end up at the points \((0, j), j = 0, 1, ..., n\), as illustrated in figure 7(a).
6.1. Partition function

The partition function $Z(n)$ for the above paths is given by the Gessel–Viennot formula.

**Theorem 6.1.**

$$Z(n) = \det_{i,j \in [0,n]} \left( \binom{2i+j}{j} \right) = 2^{n(n+1)/2}. \quad (6.1)$$

**Proof.** Let $A$ be the matrix with entries $A_{ij} = \binom{2i+j}{j}$, $i,j = 0, 1, \ldots, n$. We prove the result by $LU$ decomposition. Consider the lower triangular matrix $L$ with entries $L_{ij} = \binom{i}{j}$ and its inverse $L^{-1}$ with entries $(L^{-1})_{ij} = (-1)^{i+j} \binom{i}{j}$. Then we have the following result

$$(L^{-1}A)_{ij} = U, \quad \text{where } U_{ij} = 0 \text{ for } i > j \text{ and } U_{i,i} = 2^i.$$

Let us indeed compute:

$$(L^{-1}A)_{ij} = \sum_{k=0}^{i} (-1)^{i+k} \binom{i}{k} \binom{2k+j}{j} = \oint \frac{dt}{2i\pi t^{1+j+1}} \left( \frac{1}{t^2} - 1 \right)^i \frac{1}{(1-t)^{j+1}}$$

where the last contour integral picks up the coefficient of $t^j$. If $0 \leq j < i$ the integrand is a polynomial of degree $\leq 2i - 1$ hence the result is 0. If $j = i$, the residue at the pole $t = 1$ gives the result $2^i$. We conclude that $Z(n) = \prod_{i=0}^{n} U_{i,i} = 2^{n(n+1)/2}$ and the theorem follows. \qed

**Remark 6.2.** Note that the partition for this path model (6.1) coincides with the number of domino tilings of the Aztec diamond (3.5). There does not seem to exist any obvious natural bijection between the two sets of corresponding non-intersecting paths, and we leave this as an open question for the reader. As an indication of the difficulty of the question, we shall find that the arctic curve for the present problem is a portion of parabola, as opposed to a circle for the Aztec diamond case.

6.2. Partition functions for an escaping path

As before, we consider now the situation where the topmost endpoint $(0,n)$ is replaced by a sliding point $(0,p)$ with $p \geq n$ (see figure 7(b)). The total partition function, which we denote by $N(n,p)$, then splits as

$$N(n,p) = \sum_{\ell=0}^{2n} Z(n;\ell) W(n,p;\ell),$$

where $Z(n;\ell)$ is the partition for the same paths (bottom rectangle in figure 7(b)), but endpoints $(0,j)$, $j = 0, 1, \ldots, n-1$ and $(\ell,n)$, $\ell = 0, 1, 2, \ldots, 2n$, and $W(n,p;\ell)$ is the partition function for a single path (top rectangle in figure 7(b)) from $(\ell,n)$ to $(0,p)$, with its first step being up (to indicate that the point $(\ell,n)$ is the point when the top path exits the rectangle $[0,2n] \times [0,n]$). The latter is easily computed to be
Let us now compute the one-point function:

\[ H(n; \ell) := \frac{Z(n; \ell)}{Z(n)} = \frac{1}{2^{n^2}} \sum_{k=0}^{\min(n,2n-\ell)} \binom{n}{k}. \]

We have the following.

**Theorem 6.3.** The one-point function \( H(n; \ell) \) reads:

\[ H(n; \ell) = \frac{\tilde{U}_{n,n}}{U_{n,n}} = \frac{1}{2^n} \sum_{k=0}^{\min(n,2n-\ell)} \binom{n}{k}. \]

**Proof.** We note again that \( L^{-1} \tilde{A} = \tilde{U} \) is still upper triangular, and that

\[
\tilde{U}_{n,n} = \sum_{k} (L^{-1})_{n,k} \tilde{A}_{k,n} = \sum_{k=0}^{n} (-1)^{k+1} \binom{n}{k} \binom{n+2k-\ell}{n} \\
= \sum_{k=0}^{n} \left( \frac{1}{t^2} - 1 \right)^n \binom{n}{k} \frac{t^\ell}{(1-t)^{n+1}} \left| \frac{1-t^2}{t^\ell} \right|^{n+1} \\
= \frac{(1+t)^n}{1-t^{1+\ell}} \sum_{k=0}^{\min(n,2n-\ell)} \binom{n}{k} \\
\]

where in the second line we have used lemma 2.3. The theorem follows. \( \square \)
6.3. Tangent method

Let us evaluate the function \( N(n, p) = \sum_{\ell=0}^{2n} H(n; \ell) W(n, p; \ell) \) in the scaling limit where \( n \) is large and \( p = zn, \ell = \xi n, k = xn \). We have

\[
N(n, p) \sim \int_0^2 d\xi \int_0^{\min(1, 2-\xi)} dx \ e^{S(z, x, \xi)}
\]

where \( S(z, x, \xi) = S_0(x, \xi) + S_1(z, \xi) \), and

\[
H(n; \xi n) \sim \int_0^{\min(1, 2-\xi)} dx \ e^{-n(\log(x) + (1-x) \log(1-x) + \log(2))} \sim \int_0^{\min(1, 2-\xi)} dx \ e^{nS_0(x, \xi)}
\]

while

\[
W(n, zn; \xi n) \sim e^{n(\log(z-1+\xi) - \log(\xi-1) - \log(\xi-1))} \sim e^{nS_1(z, \xi)}.
\]

The first saddle point equation, w.r.t. \( x \) reads as follows. The extremum of \( S_0(x, \xi) \) is reached at \( x = 1/2 \) only if \( 2 - \xi > 1/2 \), i.e. \( \xi < 3/2 \), in which case the critical value of the action is \( S_0^* = 0 \). When \( \xi > 3/2 \), the extremum is at \( x = 2 - \xi \), and the critical value of the action is \( S_0^* = (2 - \xi) \log(2 - \xi) + (\xi - 1) \log(\xi - 1) + \log(2) \). The second saddle-point equation, w.r.t. \( \xi \) reads as follows. If \( \xi < 3/2 \) we have the equation:

\[
\partial_{\xi} S_1(z, \xi) = 0 \iff z - 1 + \xi = \xi
\]

which has no solution for \( z > 1 \). Therefore we have only solutions when \( \xi > 3/2 \), that obey the equation:

\[
\partial_{\xi} (S_0^* + S_1) = 0 \iff (z - 1 + \xi)(\xi - 1) = \xi(2 - \xi).
\]

This gives the relation

\[
2\xi^2 + (z - 4)\xi - (z - 1) = 0 \Rightarrow z = z(\xi) := \frac{(1 - \xi)^2}{2 - \xi}.
\]

The line through the points \((\ell, n)\) and \((0, p)\) becomes in the scaling limit the line through the points \((\xi, 1)\) and \((0, z)\), with equation

\[
y + \frac{z(\xi) - 1}{\xi} x - z(\xi) = 0 \Rightarrow 2 + 3x - 2y - (4 + 2x - y)\xi + 2\xi^2 = 0.
\]

The envelope is obtained by eliminating \( \xi \) between the above and its derivative w.r.t. \( \xi \):

\[
6 + x - 8\xi + 2\xi^2 = 0.
\]

This finally yields the parabola:

\[
-8x + 4x^2 + 8y - 4xy + y^2 = 0. \tag{6.2}
\]

As before we have derived only the portion of this curve corresponding to \( 3/2 < \xi < 2 \), i.e. \( 3/2 < x < 2 \). However, similar considerations as above show that the entire range \( 1 < \xi < 2 \), i.e. \( 0 < x < 2 \) is valid, as detailed in the next sections.
6.4. Tiling formulation and enumeration

The path model of figure 7(a) has an alternative interpretation as a rhombus tiling problem, as illustrated in figure 8(a), where the path edges cross two of the three types of tiles transversally. This gives rise to an alternative description in terms of paths across two different types of tiles. These paths have steps (1, 0) and (1, 1) and go from the points \((2i + 1/2, 0), i = 0, 1, ..., n\) to the points \((2n + 1/2, j), j = 0, 1, ..., n\). Shifting all coordinates by \((-1/2, 0)\), these are paths with the abovementioned steps from the points \((2i, 0), i = 0, 1, ..., n\) to the points \((2n, j), j = 0, 1, ..., n\). Note that this is not equivalent to a reflection of the original path problem, as the steps are different.

The partition function \(Z_2(n)\) for this model is given by the Gessel–Viennot formula.

**Theorem 6.4.** The partition function \(Z_2(n)\) reads:

\[
Z_2(n) = \det_{i,j \in [0,n]} \binom{2(n-j)}{n-j} = \det_{i,j \in [0,n]} \binom{2i}{j} = 2^{n(n+1)/2}.
\]

**Proof.** This is again proved by \(LU\) decomposition. Define the matrix \(A\) with entries \(A_{i,j} = \binom{2i}{j}\). Then the usual lower triangular matrix \(L\) with entries \(L_{i,j} = \binom{i}{j}\) is such that \(L^{-1}A = U\) is upper triangular, with diagonal elements \(U_{i,i} = 2^i\), and the theorem follows. \(\square\)

6.5. Tangent method

We now apply the tangent method, by allowing the topmost path to end at some arbitrary point \((n, p)\), thus splitting the corresponding partition function \(N_2(n, p)\) into two terms, according to the position \((\ell, n)\) of the escape point from the rectangle \([0, 2n] \times [0, n]\), where \(\ell \in [0, 2n]\). We have

\[
N_2(n, p) = \sum_{\ell=0}^{2n} Z_2(n; \ell) W_2(n, p; \ell)
\]

where \(Z_2(n; \ell)\) is the partition function for the same paths, but with the endpoint \((2n, n)\) moved to \((\ell, n)\), namely:

\[
Z_{n, \ell} = \det(\tilde{A}), \tilde{A}_{i,j} = \begin{cases} A_{i,j} & \text{if } 0 \leq j < n \\ \binom{\ell+2j-2n}{n} & \text{if } j = n \end{cases}
\]

and \(W_2(n, p; \ell)\) is the partition function of a single path from \((\ell, n)\) to \((2n, p)\), with first step \((1, 1)\) to guarantee the escape from the rectangle \([0, 2n] \times [0, n]\). We have

\[
W_2(n, p; \ell) = \binom{2n - \ell - 1}{p - n - 1}.
\]

As before we will apply the scaling analysis to the ratio \(N_2(n, p)/Z_2(n) = \sum_{\ell} H_2(n; \ell) W_2(n, p; \ell)\) where the one-point function is defined as:

\[
H_2(n; \ell) = \frac{Z_2(n; \ell)}{Z_2(n)}
\]

We have the following.
Theorem 6.5. The one-point function $H_2(n; \ell)$ reads:

$$H_2(n; \ell) = \sum_{k=0}^{\ell-n} \binom{n}{k} \left( \frac{\ell - 2k}{n} \right).$$

Proof. From the above definition of $\tilde{A}$, we see that $H_2(n; \ell) = 0$ for $\ell < n$. Using the same $L$ as before, we find that $L^{-1} \tilde{A} = \tilde{U}$, with

$$\tilde{U}_{n,n} = \sum_{k=0}^{n} (L^{-1})_{n,k} \tilde{A}_{k,n} = \sum_{k=0}^{\frac{\ell-n}{2}} (-1)^{n+k} \binom{n}{k} \left( \frac{\ell + 2k - 2n}{n} \right) = \sum_{k=0}^{\frac{\ell-n}{2}} (-1)^{k} \binom{n}{k} \left( \frac{\ell - 2k}{n} \right),$$

where we have performed a change of variables $x \to x^2$ in the first factor, and used lemma 2.3. The theorem follows.

Taking the scaling limit: $n$ large, $p = zn$, $\ell = \xi n$, $k = xn$, we find that

$$H_2(n; \xi n) \sim \int_0^{\xi-1} dx \ e^{S_0(x, \xi)}, \quad S_0(x, \xi) = -(x \log(x) + (1-x) \log(1-x))$$

and

$$W_2(n; zn; \xi \ell) \sim e^{S_1(z, \xi)}, \quad S_1(z, \xi) = (2 - \xi) \log(2 - \xi) - (z - 1) \log(z - 1) - (3 - \xi - z) \log(3 - \xi - z).$$
As before, we first extremize over $x$, with a saddle point at $x = x^* = 1/2$. Two cases must be considered (1) $\xi > 1/2$, i.e. $\xi > 3/2$, then $S_0^x = \log(2)$. (2) $\xi < 3/2$, then $x^* = \xi - 1$ and $S_0^x(\xi) = -((\xi - 1)\log(\xi - 1) - (2 - \xi)\log(2 - \xi))$.

Next we extremize over $\xi$. (1) if $\xi > 3/2$ we find the saddle point equation $(3 - \xi - z) = (2 - \xi)$ which has no solution for $z > 1$. (2) hence $\xi < 3/2$, and we have the saddle-point equation: $\partial_\xi S_0^x(\xi) + \partial_\xi S_1(z, \xi) = 0$, namely:

$$3 - \xi - z = \xi - 1 \iff \xi = 2 - \frac{z}{2}.$$  

The line through $(\xi, 1)$ and $(2, z)$ has the equation:

$$y = \frac{z - 1}{2 - \xi} - \frac{2 - z\xi}{2 - \xi} = y = \frac{z - 1}{z} - x - \frac{2 - 2z + \frac{z^2}{2}}{z} = 0$$

and the envelope is obtained by eliminating $z$ from this equation and its derivative w.r.t. $z$, namely:

$$x = 2 - \frac{z^2}{2} \quad \text{and} \quad y = z(2 - z) \quad (z \in [1, 2]).$$

Note that the limiting case $z = 2$ corresponds to the tangent at the origin, with slope 1. As expected we recover the arctic parabola equation (6.2):

$$-8x + 4x^2 + 8y - 4xy + y^2 = 0,$$

now also valid for $0 \leq x \leq 3/2$.

To summarize the methods employed above, we have represented in figure 9 the continuum limit of the arctic parabola and two sample tangents obtained by the extremization procedures above.

### 7. An interacting path model: vertically symmetric alternating sign matrices

In this last section we address the question of interacting lattice paths with the example of vertically symmetric alternating sign matrices (VSASM). It turns out that VSASM, just like ordinary ASM are in bijection with certain non-intersecting lattice paths called osculating paths, which may have ‘kissing points’ where two paths share a vertex without sharing edges, and without crossing. Assuming the applicability of the tangent method, we will derive the arctic curve for large VSASM.

#### 7.1.Vertically symmetric alternating sign matrices

An alternating sign matrix (ASM) of size $n$ is an $n \times n$ matrix satisfying the following conditions:

- All matrix elements are equal to 1, 0, or $-1$.
- The non-zero entries alternate in sign along each row and column.
- The sum of all matrix elements in any row or column is equal to 1.

Note that these three conditions taken together imply that in the first and last row of any ASM, and also in the first and last column, all matrix elements are equal to 0 except for a single 1.

A vertically symmetric alternating sign matrix (VSASM) is an alternating sign matrix of size $2n + 1$ which is also invariant under reflection about its middle column. We note that the
properties of an ASM combined with the vertically symmetric property of VSASMs implies that the central column of any VSASM has entries which alternate as $1, -1, 1, 1, -1, 1$. This implies that the unique 1 in the top and bottom row of any VSASM is located in the middle entry.

We now collect several theorems on the enumeration and refined enumeration of ASMs and VSASMs which we make use of in the tangent method calculation for VSASMs later in this section. First we state the theorems on the number of ASMs of size $n$ and VSASMs of size $2n + 1$ for any $n \in \mathbb{N}$. For ASMs we have:

**Theorem 7.1** ([16, 19]). Let $N_{\text{ASM}}(n)$ be the number of ASMs of size $n$. Then we have

$$N_{\text{ASM}}(n) = \prod_{i=0}^{n-1} \frac{(3i + 1)!}{(n + i)!}. \quad (7.1)$$

In the VSASM case we have:

**Theorem 7.2.** ([17, 18]). Let $N_{\text{VSASM}}(2n + 1)$ be the number of VSASMs of size $2n + 1$. Then we have

$$N_{\text{VSASM}}(2n + 1) = \frac{1}{2^n} \prod_{i=1}^{n} \frac{(6i - 2)!}{(4i - 1)!}. \quad (7.2)$$

Next we consider the refined enumerations of ASMs and VSASMs. For ASMs recall that the top row has all entries equal to 0 except for a single entry which is equal to 1. We may therefore consider a refined enumeration in which we compute the number of ASMs of size $n$ which have the unique 1 in their top row in the $\ell$th position, for $\ell \in [1, n]$. Then we have the following theorem.

**Theorem 7.3** ([20]). Let $N_{\text{ASM}}(n, \ell)$ be the number of ASMs of size $n$ with the unique 1 in the first row located in column $\ell$. Then we have

![Figure 9](image-url)  

*Figure 9.* The arctic parabola in the continuum limit separates frozen phases (top left and right corners) from a disordered phase (bottom). We have represented two sample tangents obtained from the methods above, together with their intercepts on the line $y = 1$. 
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\[ N_{\text{ASM}}(n, \ell) = \frac{(n + \ell - 2)}{(n - 1)} \left( \frac{2n - 1 - \ell}{n - 1} \right) \frac{3n - 2}{n - 1} N_{\text{ASM}}(n). \]  

(7.3)

Note also that by the 90° rotation symmetry of the ASM enumeration problem, \( N_{\text{ASM}}(n, \ell) \) is also equal to the number of ASMs of size with their unique 1 at position \( \ell \) in the first column, or in the last column, or in the last row.

In the case of VSASMs we already remarked that the unique 1 in the first and last rows must appear in the central column. However, the unique 1 in the first column can still take any position (and likewise for the last column). We can therefore consider a refined enumeration in which we compute the number of VSASMs of size 2\( n + 1 \) in which the unique 1 in the first column is at position \( \ell \) for \( \ell \in [1, 2n + 1] \). Then we have:

**Theorem 7.4 ([18]).** Let \( N_{\text{VSASM}}(2n + 1, \ell) \) be the number of VSASMs of size 2\( n + 1 \) with the unique 1 in the first column located in row \( \ell \). Then we have

\[ N_{\text{VSASM}}(2n + 1, \ell) = \frac{N_{\text{VSASM}}(2n - 1)}{(4n - 2)!} \sum_{i=1}^{\ell-1} (-1)^{\ell+i-1} \frac{(2n + i - 2)!(4n - i - 1)!}{(i - 1)!(2n - i)!}. \]  

(7.4)

Due the reflection symmetry of VSASMs, \( N_{\text{VSASM}}(2n + 1, \ell) \) is also the number of VSASMs of size 2\( n + 1 \) with the unique 1 in the last column located in row \( \ell \). As mentioned before, such an alternating formula is not suitable for large size estimates, as terms may cancel out. Fortunately, we will be able to use the following formula relating the refined enumerations of ASMs and VSASMs. More precisely, the formula relates certain generating functions formed from the refined enumerations of ASMs and VSASMs. We have:

**Theorem 7.5 ([18]).** The following equality holds:

\[ \frac{1}{N_{\text{VSASM}}(2n - 1)} \sum_{\ell=1}^{2n} N_{\text{VSASM}}(2n + 1, \ell) t^{\ell-1} = \frac{1}{N_{\text{ASM}}(2n - 1)} t + \frac{1}{t + 1} \sum_{\ell=1}^{2n} N_{\text{ASM}}(2n, \ell) t^{\ell-1}. \]  

(7.5)

Note also that on the left-hand side of this equation we can extend the range of the sum to \( \ell \in [1, 2n + 1] \) due the fact that \( N_{\text{VSASM}}(2n + 1, 1) = N_{\text{VSASM}}(2n + 1, 2n + 1) = 0 \) by the properties of VSASMs.

### 7.2. Mapping to six vertex model configurations

The ASMs of size \( n \) are in bijection with configurations of the six vertex model with *domain wall* boundary conditions (DWBCs) on an \( n \times n \) portion of the square lattice \( \mathbb{Z}^2 \). We now give a short review of the six vertex model and the DWBCs, and then discuss the mapping from configurations of the six vertex model with DWBCs to ASMs.

The six vertex model is a 2D statistical mechanical model in which the degrees of freedom are arrows placed on the edges of the square lattice \( \mathbb{Z}^2 \). In addition, the arrow configurations are constrained to satisfy the *ice rule*, which states that at each vertex the number of incoming arrows must be equal to the number of outgoing arrows. On the square lattice there are exactly six possible arrow configurations around a vertex that satisfy the ice rule, thus explaining the
name of this model. The six possible vertex configurations are traditionally divided into three types called $a$, $b$, and $c$-type vertices. Finally, there is also a representation of configurations of the six vertex model in terms of lattice paths. In this representation each possible configuration of arrows around a vertex is mapped to a segment of one or two lattice paths. We display the six possible vertex configurations of this model, along with their type ($a$, $b$, or $c$) and their representation in terms of lattice paths in figure 10. The additional data in the last row of figure 10 (namely, the numbers 0, 1 or $-1$) is related to the mapping to ASMs and will be explained shortly.

We note here that in the lattice path formulation the six vertex model should be thought of as a model of interacting lattice paths due to the fact that two paths are allowed to touch, but not cross, at a single vertex (see the first $a$-type vertex in figure 10). This should be contrasted with the other lattice path models studied in this paper in which the paths are all non-intersecting. In non-intersecting lattice path models two paths are never allowed to meet at a single vertex. We may therefore characterize the non-intersecting lattice path models as models of non-interacting lattice paths.

We now describe the DWBCs for the six vertex model, which were first considered by Korepin [14]. The six vertex model with DWBCs simply consists of the six vertex model formulated on an $n \times n$ region of the square lattice, in which the arrows on the edges which ‘stick out’ from the boundary are fixed to point in towards the domain on the left and right boundaries, and out of the domain on the top and bottom boundaries. In this lattice path formulation the DWBCs imply that on an $n \times n$ domain there are $n$ lattice paths which enter at the top of the domain and exit on the left side of the domain. A sample configuration of the six vertex model with DWBCs, in both the arrow and lattice paths formulations, is shown for a system of size $n = 5$ in figure 11.

To define the partition function for the six vertex model we assign weights to each of the six possible vertex configurations. We will make a slight abuse of notation and use the letters $a$, $b$, and $c$ to also denote the weights of the vertices of types $a$, $b$, and $c$, respectively. The weight $w(\mathcal{C})$ of a configuration $\mathcal{C}$ of the six vertex model is defined to be equal to the product of the weights of all vertices in the configuration $\mathcal{C}$. The partition function $Z_{6VDWBC}(n; a, b, c)$ for the six vertex model with DWBCs for a system of size $n$ is defined as

$$Z_{6VDWBC}(n; a, b, c) = \sum_{\mathcal{C}} w(\mathcal{C}), \quad (7.6)$$

where the sum is taken over all configurations $\mathcal{C}$ consistent with the DWBCs. Here we have indicated explicitly that the partition function $Z_{6VDWBC}(n; a, b, c)$ is a function of the weights $a$, $b$, and $c$.

It is known that the ASMs of size $n$ are in bijection with the configurations of the six vertex model on an $n \times n$ grid and with DWBCs. The mapping from configurations of the six vertex model to ASMs is a set of rules that convert a vertex in the six vertex model to a matrix element (1, 0, or $-1$) in an ASM. We exhibit this rule in the last row of figure 10. As an example, the six vertex model configuration shown in figure 11 maps to the following ASM of size 5:

$$\begin{pmatrix}
0 & 0 & 1 & 0 & 0 \\
1 & 0 & -1 & 0 & 1 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & -1 & 1 & 0 \\
0 & 0 & 1 & 0 & 0
\end{pmatrix} \quad (7.7)$$

We note here that this particular ASM is also a VSASM.
Since the ASMs of size \( n \) are in bijection with the configurations of the six vertex model with DWBCs on a domain of size \( n \), we have that

\[
Z_{6\text{VDWBC}}(n; 1, 1, 1) = N_{\text{ASM}}(n).
\] (7.8)

The particular parameter values \( a = b = c = 1 \) correspond to the so-called ice point. In what follows we will need a modified partition function for the six vertex model with DWBCs in which we only sum over configurations \( \tilde{C} \) which can be mapped to a VSASM. Clearly, this modified partition function only exists for a square domain of odd size. We denote this modified partition function by \( \tilde{Z}_{6\text{VDWBC}}(2n + 1; a, b, c) \), and we have

\[
\tilde{Z}_{6\text{VDWBC}}(2n + 1; a, b, c) = \sum_{\tilde{C}} w(\tilde{C}),
\] (7.9)

where again the sum is only over those configurations of the six vertex model with DWBCs on a \((2n + 1) \times (2n + 1)\) grid which map to a VSASM. From the definition of this modified partition function it is also clear that

\[
\tilde{Z}_{6\text{VDWBC}}(2n + 1; 1, 1, 1) = N_{\text{VSASM}}(2n + 1).
\] (7.10)

As a side note, we remark here that the usual approach to compute \( N_{\text{VSASM}}(2n + 1) \) and the refined enumeration \( N_{\text{VSASM}}(2n + 1, \ell) \) uses a mapping from VSASMs to configurations of the six vertex model with ‘U-turn’ boundary conditions, as originally considered by Kuperberg [17]. We do not follow this approach here since the only ingredient we will need for the application to the tangent method is theorem 7.5 relating the refined enumeration of VSASMs to that of ASMs. We now move on to the discussion of the tangent method for this system.

7.3. Tangent method

We now apply the tangent method to compute the arctic curve for VSASMs. We will prove that the arctic curve for VSASMs is actually identical to the arctic curve of ordinary ASMs. Any possible differences between these two curves disappear in the scaling limit. We should mention that in the context of VSASMs, the arctic curve can be understood as separating a
‘disordered’ interior region of a VSASM, in which 1, 0, and −1 entries occur, and an ordered or frozen region in which only zeros appear. For VSASMs (and also for ASMs) the arctic curve will actually touch each of the four boundaries of the domain at a single point, and this point corresponds to the most likely location for the unique 1 entry in the last row or column of the VSASM (or ASM).

To apply the tangent method we need to extend the original domain of the six vertex model with DWBCs. In their paper, Colomo and Sportiello chose to extend the domain vertically. Here we choose to extend the domain horizontally instead, as this type of extension is more natural for VSASMs. The reason is that in the case of VSASMs, the path (in the six vertex model formulation) which starts at the bottom left of the domain is constrained to turn upwards when it reaches the center of the domain, and not before or after. Therefore it would not make sense to extend the domain vertically by allowing this last path to dip below the original domain, since for VSASMs it would always be constrained to pass from the original domain into the extended portion at the central column. On the other hand, this last path can arrive at the far right column of the original domain at any height \( \ell > 1 \), so it does make sense to extend the domain horizontally and allow this last path to pass into the extended portion of the domain after hitting the far right column of the original domain at the height \( \ell \). This is the approach that we take in this section.

To extend the domain we introduce a new integer \( k > 0 \) and extend the domain to a region of the square lattice of size \((2n + 1 + k) \times (2n + 1)\). We label the vertices in the domain by Cartesian coordinates \((x, y) = (i, j)\) with \( i = 1, \ldots, 2n + 1 + k \) and \( j = 1, \ldots, 2n + 1 \), and we consider \( 2n + 1 \) lattice paths on this domain. All \( 2n + 1 \) lattice paths enter at the left side of the domain. The top-most \( 2n \) lattice paths exit at the top of the domain at \( x \) positions \( i = 1, \ldots, 2n \), while the final lattice path exits the domain at the top of the very last column with coordinate \( i = 2n + 1 + k \). In addition, we require that the portion of the lattice path configurations which lies is the original domain is such that it would map to a VSASM under the mapping from lattice paths/six vertex model configurations to ASMs (more precisely, it should map to a VSASM if we turn the last path upwards at \( i = 2n + 1 \) so that it exits in the last column of the original domain). An example of an allowed lattice path configuration on this extended domain is shown in figure 12. It corresponds to the example shown in figure 11, but where we let the last path cross into the extended part of the domain.

After rescaling all coordinates by a factor of \( 2n + 1 \), the original six vertex model domain is rescaled to lie in the unit square \( x \in [0, 1], y \in [0, 1] \). Due the way we have chosen to extend the domain, the tangent method will naturally produce the portion of the arctic curve
for VSASMs which lies in the quadrant \( x \in [\frac{1}{2}, 1], y \in [0, \frac{1}{2}] \), and we will have to keep this in mind when comparing our final expression for the arctic curve with other expressions in the literature for the arctic curve of ASMs. Since any given VSASM has a ‘mirror image’ obtained by reflecting this VSASM about the central row, we can argue by symmetry that the portion of the arctic curve in the region \( x \in [\frac{1}{2}, 1], y \in [\frac{1}{2}, 1] \) can be obtained from the curve in the region \( x \in [\frac{1}{2}, 1], y \in [0, \frac{1}{2}] \) by simply reflecting it over the line \( y = \frac{1}{2} \). Finally, due to the defining symmetry of VSASMs (which is the symmetry of reflection about the central column), the entire arctic curve in the region with \( x \leq \frac{1}{2} \) is obtained from the arctic curve in the region \( x \geq \frac{1}{2} \) by reflection over the line \( x = \frac{1}{2} \).

As in previous sections, we compute the full partition function for the problem on the extended domain by expanding in the vertical position \( \ell \) at which the last path crosses from the original domain into the extended domain. In what follows we use the notation \( \tilde{n} := 2n + 1 \) for convenience. If we denote by \( N(\tilde{n}, \ell) \) the partition function for the six vertex model on the extended domain, then we can write

\[
N(\tilde{n}, \ell) = \sum_{\ell=1}^{\tilde{n}} Z(\tilde{n}, \ell) Y(\ell, k) .
\] (7.11)

Here \( Z(\tilde{n}, \ell) \) is the partition function for paths on the \( \tilde{n} \times \tilde{n} \) portion of the domain in which the last path exits this portion of the domain at the height \( \ell \). The other factor \( Y(\ell, k) \) is the partition function for the piece of the last path which lies outside of the original \( \tilde{n} \times \tilde{n} \) portion of the domain. For the application to VSASMs we only need to consider the six vertex model at the ice point \( a = b = c = 1 \), and so we restrict our attention to this case in what follows.

Since we work at the ice point, we find that \( Z(\tilde{n}, \ell) \) is exactly equal to the number of VSASMs of size \( \tilde{n} \) whose unique 1 in the last column is at position \( \ell \),

\[
Z(\tilde{n}, \ell) = N_{\text{VSASM}}(\tilde{n}, \ell) .
\] (7.12)

Next, \( Y(\ell, k) \) is equal to the total number of six vertex model paths which start at \((\tilde{n}, \ell)\), end at \((\tilde{n} + k, \tilde{n} + 1)\), and are constrained to have the first step be horizontal and the last step be vertical. This constraint then implies that \( Y(\ell, k) \) is just equal to the total number of six vertex model paths which start at \((\tilde{n} + 1, \ell)\) and end at \((\tilde{n} + k, \tilde{n})\). By a ‘six vertex model path’ we mean a path whose segments are constrained to be of the type illustrated in figure 10. In addition, if we have only a single path, then we are restricted to the vertices of type \( b \) and \( c \) from figure 10. This calculation is presented in, for example, appendix 1 of [4], and the answer can be expressed as

\[
Y(\ell, k) = \sum_{p=0}^{\min(k-1, \tilde{n}-\ell)} \binom{k-1}{p} \binom{\tilde{n}-\ell}{p} .
\] (7.13)

Here the summation variable \( p \) can be interpreted as counting the number of ‘north-east corners’ of a path, where we have a \((0, 1)\) step immediately followed by a \((1, 0)\) step (alternatively, the second type \( c \) vertex from figure 10). Thus, we find that the total partition function of the six vertex model (with our restriction to configurations in the \( \tilde{n} \times \tilde{n} \) portion such that they map to a VSASM) at the ice point on the extended domain is

\[
N(\tilde{n}, \ell) = \sum_{\ell=1}^{\tilde{n}} \sum_{p=0}^{\min(k-1, \tilde{n}-\ell)} N_{\text{VSASM}}(\tilde{n}, \ell) \binom{k-1}{p} \binom{\tilde{n}-\ell}{p} .
\] (7.14)

For later use we also define the one-point function \( H(\tilde{n}, \ell) \) as
To compute the asymptotics of this partition function we first introduce new scaled variables \( z, \xi, \) and \( \eta \) as \( k = \tilde{n}z,\ell = \tilde{n}\xi,\) and \( p = \tilde{n}\eta. \) We then follow Colomo and Sportiello and define the ‘free energy’

\[
F[z] = \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \log \left[ \frac{N(\tilde{n},k)}{N_{\text{VSASM}}(\tilde{n})} \right]
\]

and the ‘action’

\[
S(\xi,\eta;z) = \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \log \left[ \binom{k-1}{p} \binom{\tilde{n}-\ell}{p} H(\tilde{n},\ell) \right].
\]

After rescaling of the variables and applying Stirling’s formula \( \log(q!) \approx q \log(q) - q \) for large \( q, \) the action takes the form

\[
S(\xi,\eta;z) = f(z) - 2f(\eta) - f(z - \eta) + f(1 - \xi) - f(1 - \xi - \eta) + \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \log [H(\tilde{n},\xi\tilde{n})],
\]

where we used the abbreviated notation \( f(x) := x \log(x) - x. \) In the scaling limit the free energy is dominated by a contribution from the stationary point of the action, which can be found by solving the equations

\[
\frac{\partial S(\xi,\eta;z)}{\partial \xi} = 0 \quad \text{(7.19)}
\]

\[
\frac{\partial S(\xi,\eta;z)}{\partial \eta} = 0 \quad \text{(7.20)}
\]

for the saddle point solutions \( \xi_{sp}(z) \) and \( \eta_{sp}(z) \) which are functions of \( z. \)

Before rescaling of the coordinates, the tangent line we are looking for starts at the point \((\tilde{n},\ell)\) and ends at \((\tilde{n} + k,\tilde{n}).\) After rescaling and solving for \( \xi_{sp}(z), \) the most likely value of the rescaled \( y \) coordinate \( \ell \) where the last path leaves the original domain, this line is defined by the equation

\[
H(\tilde{n},\ell) = \frac{N_{\text{VSASM}}(\tilde{n},\ell)}{N_{\text{VSASM}}(\tilde{n})}.
\]
\[ y = \left( \frac{1 - \xi_{sp}(z)}{z} \right) x + \xi_{sp}(z) - \left( \frac{1 - \xi_{sp}(z)}{z} \right). \] (7.21)

We now proceed with the solution of the saddle point equations. After differentiating with respect to \( \xi \) and \( \eta \), the two saddle point equations take the form

\[ \text{Log} \left( \frac{1 - \xi - \eta}{1 - \xi} \right) + \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \frac{\partial}{\partial \xi} \text{Log} \left[ H(\tilde{n}, \xi_{\tilde{n}}) \right] = 0 \] (7.22)

and

\[ \text{Log}(z - \eta) - 2\text{Log}(\eta) + \text{Log}(1 - \xi - \eta) = 0. \] (7.23)

The second equation can be solved immediately and allows us to express \( \eta_{sp}(z) \) in terms of \( \xi_{sp}(z) \) as

\[ \eta_{sp}(z) = \left( \frac{1 - \xi_{sp}(z)}{1 - \xi_{sp}(z) + z} \right) z. \] (7.24)

Solving the first saddle point equation is more complicated due to the presence of the term involving the one point function \( H(\tilde{n}, \xi_{\tilde{n}}) \). Before presenting the solution we first eliminate \( \eta \) from this equation using the solution of the second saddle point equation to find

\[ \text{Log} \left( \frac{1 - \xi}{1 - \xi + z} \right) + \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \frac{\partial}{\partial \xi} \text{Log} \left[ H(\tilde{n}, \xi_{\tilde{n}}) \right] = 0. \] (7.25)

To solve this equation we follow the method of [4] and find that the solution \( \xi_{sp}(z) \) is given implicitly by the equation

\[ \xi_{sp}(z) = r(t), \] (7.26)

obtained by inverting the relation:

\[ t = \frac{1 - \xi_{sp}(z)}{1 - \xi_{sp}(z) + z}, \] (7.27)

with

\[ r(t) = \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \frac{d}{dt} \text{Log}(h_{\tilde{n}}(t)), \] (7.28)

and where we introduced the generating function

\[ h_{\tilde{n}}(t) := \sum_{\ell=1}^{\tilde{n}} H(\tilde{n}, \ell) t^{\ell-1} \] (7.29)

for the one point function \( H(\tilde{n}, \ell) \). To obtain equations (7.26) and (7.27), we simply estimate for large \( \tilde{n} \):

\[ h_{\tilde{n}}(t) \sim \int_{0}^{t} d\xi \; H(\tilde{n}, \xi_{\tilde{n}}) e^{\xi S(\xi, t)} = \int_{0}^{1} d\xi \; e^{\tilde{n} S(\xi, t)}, \]

with

\[ S(\xi, t) = \xi \text{Log}(t) + \frac{1}{\tilde{n}} \text{Log} \left[ H(\tilde{n}, \xi_{\tilde{n}}) \right]. \]
If we compare the saddle-point equation \( \partial_\xi S(\xi, t) = 0 \) with equation (7.25), then we find a relation between \( \xi_{sp}(z) \) and \( t \), namely (7.27). Finally, the large \( \tilde{n} \) estimate of \( h_\tilde{n}(t) \) is given by:

\[
h_\tilde{n}(t) \sim e^{S(\xi_{sp}(z); t)}
\]

therefore

\[
r(t) := \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \frac{d}{dt} \log(h_\tilde{n}(t)) = \lim_{\tilde{n} \to \infty} \left( \xi + \frac{t}{\tilde{n}} \frac{d}{dt} \log(h_\tilde{n} + \frac{1}{\tilde{n}} \frac{d}{dt} \log[H(\tilde{n}, \xi)]) \right) \bigg|_{\xi = \xi_{sp}(z)}
\]

\[
= \xi_{sp}(z) + t \frac{d}{dt} \log([H(\tilde{n}, \xi)]) \bigg|_{\xi = \xi_{sp}(z)} = \xi_{sp}(z)
\]

by use of the saddle point equation (7.25) and the relation (7.27). This gives (7.26).

We now note that due to the relation

\[
\frac{1 - \xi_{sp}(z)}{z} = \frac{t}{1 - t},
\]

and the fact that \( \xi_{sp}(z) = r(t) \), we can eliminate \( \xi_{sp}(z) \) and \( z \) in favor of \( t \) and \( r(t) \) in our expression for the tangent line. After this change of variables the equation for the tangent line takes the form

\[
y = \left( \frac{t}{1 - t} \right) x + r(t) - \frac{t}{1 - t}.
\]

To solve for the arctic curve we define the function

\[
F(x, y; t) = -y + \left( \frac{t}{1 - t} \right) x + r(t) - \frac{t}{1 - t}
\]

such that the equation \( F(x, y; t) = 0 \) is the equation defining the tangent line. To extract the arctic curve we then solve the simultaneous equations \( F(x, y; t) = 0 \) and \( \frac{\partial}{\partial y} F(x, y; t) = 0 \) to solve for \( x \) and \( y \) in terms of \( t \), which then yields a parametric form \((x(t), y(t))\) for the arctic curve.

Our final task is to solve for the function \( r(t) \). To start, we write out the generating function \( h_\tilde{n}(t) \) in more detail (using the definition of the one point function) as

\[
h_\tilde{n}(t) = \sum_{\ell=1}^{\tilde{n}} \frac{N_{VSASM}(\tilde{n}, \ell)}{N_{VSASM}(\tilde{n})} t^{\ell-1}.
\]

Next we apply theorem 7.5 to find that (recall that \( \tilde{n} = 2n + 1 \))

\[
h_\tilde{n}(t) = \frac{N_{VSASM}(\tilde{n} - 2)}{N_{VSASM}(\tilde{n}) N_{ASM}(\tilde{n} - 2)} \frac{t}{t + 1} \sum_{\ell=1}^{\tilde{n} - 1} N_{ASM}(\tilde{n} - 1, \ell) t^{\ell-1} = \frac{N_{VSASM}(\tilde{n} - 2) N_{ASM}(\tilde{n} - 1)}{N_{VSASM}(\tilde{n}) N_{ASM}(\tilde{n} - 2)} \frac{t}{t + 1} \sum_{\ell=1}^{\tilde{n} - 1} N_{ASM}(\tilde{n} - 1, \ell) t^{\ell-1} = \frac{N_{VSASM}(\tilde{n} - 2) N_{ASM}(\tilde{n} - 1)}{N_{VSASM}(\tilde{n}) N_{ASM}(\tilde{n} - 2)} \frac{t}{t + 1} h_{\tilde{n} - 1}(t),
\]

(7.34)
where we defined

$$h_{\tilde{n}-1}(t) = \sum_{\ell=1}^{\tilde{n}-1} \frac{N_{\text{ASM}}(\tilde{n}-1, \ell)}{N_{\text{ASM}}(\tilde{n}-1)} t^{\ell-1}.$$  \hspace{1cm} (7.35)

We now compute the function $r(t)$ as

$$r(t) = \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \frac{d}{dt} \log(h_{\tilde{n}}(t)) = \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \left[ \frac{1}{1+t} + t \frac{d}{dt} \log(h_{\tilde{n}}(t)) \right].$$  \hspace{1cm} (7.36)

The term $\frac{1}{\tilde{n}}$ came from differentiating the factor of $\frac{1}{\tilde{n}}$ which appeared in theorem 7.5. In the $\tilde{n} \to \infty$ limit this term goes to zero due to the prefactor of $\frac{1}{\tilde{n}}$ and, since $\tilde{n}-1$ can be replaced with $\tilde{n}$ in this limit, we find that

$$r(t) = r_{\text{ASM}}(t),$$  \hspace{1cm} (7.37)

where

$$r_{\text{ASM}}(t) = \lim_{\tilde{n} \to \infty} \frac{1}{\tilde{n}} \frac{d}{dt} \log(h_{\tilde{n}}(t)) = \frac{\sqrt{t^2 - t + 1} - 1}{t - 1}.$$  \hspace{1cm} (7.38)

To close this section we present the explicit form of the arctic curve for VSASMs which, due to the relation $r(t) = r_{\text{ASM}}(t)$, is identical to the arctic curve for ordinary ASMs. Solving the simultaneous equations $F(x, y; t) = 0$ and $\frac{\partial}{\partial t} F(x, y; t) = 0$ using the explicit form of $r_{\text{ASM}}(t)$ yields the parametric form of the arctic curve,

$$x(t) = \frac{1 + t}{2\sqrt{1 - t + t^2}}$$  \hspace{1cm} (7.40)

$$y(t) = \frac{-2 + t + 2\sqrt{1 - t + t^2}}{2\sqrt{1 - t + t^2}}.$$  \hspace{1cm} (7.41)

It is now straightforward to verify that $x(t)$ and $y(t)$ satisfy the equation

$$4(1 - x) - 4(1 - x)^2 + 4y - 4y^2 + 4(1 - x)y - 1 = 0,$$  \hspace{1cm} (7.42)

which is exactly the equation for the portion of the arctic curve for ASMs in the region $x \in [\frac{1}{2}, 1]$, $y \in [0, \frac{1}{2}]$ of the rescaled domain. Note that in [3] Colomo and Pronko obtained the portion of the arctic curve for ASMs in the region $x \in [0, \frac{1}{2}]$, $y \in [0, \frac{1}{2}]$, and so our curve differs from theirs by the replacement $x \to 1 - x$, which implements the reflection over the line $x = \frac{1}{2}$. It is also interesting and worth noting that we still obtain the correct arctic curve for ASMs and VSASMs by extending the domain of the six vertex model horizontally instead of vertically as in [4]. Finally, let us also note that due to the term $\frac{1}{\tilde{n}}$ appearing in theorem 7.5, one would naively think that the arctic curves for ASMs and VSASMs would be different. However, the contribution of this extra term to the function $r(t)$ is suppressed in the $\tilde{n} \to \infty$ limit, and so the arctic curves for ASMs and VSASMs turn out to be identical. However,
finite-size numerical studies of the arctic curve in VSASMs should show slight differences from the arctic curve for ASMs, due to this extra term.

8. Conclusion

8.1. Summary

In this paper, we have explored four specific statistical models that can be rephrased into (weakly) non-intersecting path models, and shown how to apply the tangent method for determining the arctic curve in the limit of large size. Our exact results confirm the applicability of the method in the case of the domino tilings of the Aztec diamond (section 3), where the arctic circle was derived rigorously by other methods [1, 7]. Our Dyck path model for the rhombus tiling of a half-hexagon (sections 4 and 5) and the osculating path model for VSASM (section 7) both lead to analogous results: the arctic curve for a half-domain with fixed boundary conditions along the cut is identical to that of the full domain (ellipse for the hexagon, and 4 quarters of ellipse for ASM). Finally, the other path model of section 6 leads to a parabola, which, despite the fact that the partition function matches that of domino tilings of the Aztec diamond, points to a very different asymptotic behavior, which we conjecture can be traced back to the choice of boundary conditions.

8.2. Open problems

We are left with many open questions. The tangent method itself remains to be proved rigorously. Regarding its range of applicability, it seems to not only apply to strictly non-intersecting paths, but to weakly interacting ones as well, for which kissing or osculating points are allowed. We may perhaps test this on other path models, such as osculating large Schroeder paths, corresponding to tilings of the Aztec diamond or other domains by means of $2 \times 1$ and $1 \times 2$ dominos and a finite set of extra larger tiles that account for the various kissing point configurations. Note in that case the novel possibility for three paths to share an osculating vertex. Other cases of interest are the path models corresponding to higher spin versions of the 6 Vertex model (such as the spin-1 19 vertex model for instance). Finally, another set of problems concerns path models with inhomogeneous weights, namely with steps of different weights depending on their positions. For instance periodic inhomogeneous weight domino tilings of the Aztec diamond were studied in [7], and shown to give rise to more involved arctic curves, including 'bubbles' of intermediate disorder semi-crystalline phases. We intend to return to these problems in a future publication.

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