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Machine-Learned Dark Matter Subhalo Candidates in the 4FGL-DR2: Search for the Perturber of the GD-1 Stream

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Abstract. The detection of dark matter subhalos without a stellar component in the Galactic halo remains a challenge. We use supervised machine learning to identify high-latitude gamma-ray sources with dark matter-like spectra among unassociated gamma-ray sources in the 4FGL-DR2. Out of 843 4FGL-DR2 unassociated sources at $|b| \geq 10^\circ$, we select 73 dark matter subhalo candidates. Of the 69 covered by the \textit{Neil Gehrels Swift Observatory (Swift)}, 17 show at least one X-ray source within the 95\% LAT error ellipse and 52 where we identify no new sources. This latest inventory of dark subhalos candidates allows us to investigate the possible dark matter substructure responsible for the perturbation in the GD-1 stellar stream. In particular, we examine the possibility that the alleged GD-1 dark subhalo may appear as a 4FGL-DR2 gamma-ray source from dark matter annihilation into Standard Model particles.
1 Introduction

It has been long hypothesized that the Milky Way halo is populated by thousands of dark matter subhalos above a mass of $10^6 M_\odot$ [1–3]. With dedicated optical surveys, the number of detected dwarf galaxies orbiting the Milky Way has grown significantly and sets a lower bound on the abundance of dark matter subhalos [4]. However, there might be dark substructures that are too faint or too distant to be picked up by optical surveys alone [5]. Localizing dark matter subhalos without detectable stellar populations is still of critical importance not only to understand structure formation but also to reveal the particle nature of dark matter directly.

Weakly interacting massive particles (WIMPs) in the nearest and most massive subhalos could produce observable gamma-ray sources with significantly curved energy spectra from annihilating dark matter [6]. In general, one can parametrize WIMP annihilation spectra with a super-exponential cutoff power law [7, 8]:

$$dN_{DM}/dE(E) = K \left( \frac{E}{E_0} \right)^{-\Gamma} e^{-\left( \frac{E}{E_{cut}} \right)^\beta},$$

where $K$ is a prefactor, $E_0 = 10^3$ MeV is the pivot energy, $\Gamma$ is the spectral photon index, $E_{cut}$ is the cutoff energy and $\beta$ is the curvature index. Taking advantage of the extraordinary map of the gamma-ray sky produced by the Large Area Telescope (LAT) on the Fermi Gamma-ray Space Telescope (Fermi), several searches have used source catalogs released by the LAT collaboration to look for this distinctive spectral feature among unassociated sources [7–19]. Although there are interesting subhalo candidates with exponential cutoffs, it is still not yet possible to distinguish them from gamma-ray pulsars [16].

An alternate approach is based on the idea that dark matter subhalos orbiting in the Milky Way halo might induce gaps when they cross elongated dynamically-cold stellar streams...
With astronomical surveys such as the Gaia mission, there is now sufficient precision to study stellar streams in extreme detail. The GD-1 stellar stream is currently the most notable example of a possible stream crossing [25, 26]. The observed morphology of GD-1 (gap and off-stream spur of stars) is naturally reproduced in models of the stellar stream that include an encounter with a massive object (∼10^6 – 10^8 M☉), like a globular cluster, a dwarf galaxy, or a dark matter subhalo [26, 27].

Each of these techniques offers excellent opportunities to survey the subhalo population. But a combination of approaches might boost our chances for success. In this paper, we use the first incremental version of the Fourth Fermi-LAT Catalog of gamma ray sources (4FGL-DR2, for Data Release 2) to select possible dark matter subhalos using machine-learning techniques [12, 18]. Out of 843 unassociated gamma-ray sources at |b| ≥ 10°, we extract a subset of 73 sources that are consistent with the spectral shape of annihilating dark matter. We then search for X-ray sources in the LAT 95% error ellipses using archival Swift data. We next compare the locations of the subhalo candidates with the predicted locations of the GD1’s perturbing subhalo from Gaia.

2 Dark Subhalo Selection

2.1 Dataset

We start with the 4FGL-DR2 covering the time period August 4, 2008, to August 2, 2018 in the 50 MeV-1 TeV energy range [28, 29]. For our study, we use the latest version available at the time of this writing gll-psc-v27.fit. These sources are extremely well characterized by 74 catalog columns. Of the 5787 4FGL-DR2 sources, 1670 are unassociated. This leaves 4117 associated sources. Since we are only interested in searching for dark subhalo candidates away from the Galactic plane, we can safely remove Galactic plane sources such as supernova remnants and pulsar wind nebulae. The largest source population at high Galactic latitude is extragalactic. We parametrize these sources with an extragalactic set that includes 3505 training sources from all AGN classes in the 4FGL. The extragalactic set includes 11 non-blazar active galaxies (AGN,agn) 1382 blazar candidates of uncertain type (BCU,bcu), 1308 BL Lacs (BLL,bll), 743 flat-spectrum radio quasars (FSRQ,fsrq), 9 narrow-line Seyferts 1 (NLSY1,nlsy1), 44 radio galaxies (RDG,rdg), 2 soft spectrum radio quasars (ssrq), 5 compact steep spectrum quasar (css), and 1 Seyfert galaxy (sey).

In order to generate a model for dark subhalo searches we harness the vast catalog of gamma-ray pulsar and globular cluster spectra [18]. Early studies of unassociated gamma-ray sources led to the important realization that gamma-ray pulsar spectra would be nearly indistinguishable from the super-exponential spectra generated by WIMP annihilation in dark subhalos [6]. Thus formally, a high-latitude, non-variable gamma-ray pulsar candidate without detected pulsation is also a potential subhalo candidate. For our subhalo set, we include 235 identified pulsars (PSRs), 30 globular clusters (glc), and 24 associated pulsars (psrs) with spectra that are also consistent with annihilations of ∼20–70 GeV dark matter particles [16, 17].

We end up with a grand total of 3794 training sources. With 3505 extragalactic and 289 pulsars, our training set is imbalanced. Whenever there is imbalance, machine learning classification tends to be artificially skewed towards the majority class. As in [18], we use the Synthetic Minority Over-sampling Technique to balance the class distribution (SMOTE, [30]). SMOTE adds synthetic pulsars sampled from the existing pulsars using k-nearest neighbors until they both contain near equal samples.
2.2 Machine Learning Algorithms

For 4FGL-DR2 training, testing and classification, we rely on the scheme introduced in [18]. For training we use Random Forest and eXtreme Gradient Boosting (XGBoost) tree algorithms. Although both rely on classification trees, Random Forest uses feature bagging that makes decisions by a committee of individual classification trees. Each tree casts a vote for the predicted class and the majority vote is taken as the prediction, thus it reduces variance [31]. XGBoost \(^1\) tries to do better at each interaction, thus it reduces bias [32]. We settle on a combination of R and scikit-learn \(^2\) for the rest of the paper.

2.3 Feature Selection

Given the large number of 4FGL-DR2 features, it is critical to determine their relative importance for correct classification. To do this we rely on the Gini coefficient as determined by Random Forest [33, 34]. A larger value of Gini importance means that a particular feature plays a greater role in predicting a certain class. This is shown in Figure 1. As shown, the most important predictors are:

- **LP_SigCurv:** Significance (in \(\sigma\) units) of the fit improvement between PowerLaw and LogParabola.
- **PLEC_SigCurv:** Same as LP_SigCurv for PLSuperExpCutoff model
- **LP_beta:** Curvature parameter when fitting with LogParabola
- **Flux1000:** Integral photon flux from 1 to 100 GeV
- **Frac_Variability:** Fractional variability computed from the fluxes in each year
- **Variability_Index:** Sum of \(2\times\log(\text{Likelihood})\) difference between the flux fitted in each time interval and the average flux over the full catalog interval
- **PLEC_Index:** Low-energy photon index when fitting with PLSuperExpCutoff
- **PL_Index:** Photon index when fitting with PowerLaw
- **Pivot_Energy:** Energy at which error on differential flux is minimal
- **LP_Index:** Photon index at Pivot Energy when fitting with LogParabola

Another reliable importance measure is the evaluation of “out-of-bag” (OOB) performance as a function of number of features. As can be seen Figure 1, the OOB performance tends to settle after 8 or more features are included. Since there is no obvious learning loss, we decided to retain the top 10 features for the rest of the classification.

\(^{1}\)https://github.com/dmlc/xgboost
\(^{2}\)http://www.scikit-learn.org
2.4 Training and Classification

For cross validation, we divide the dataset into training (70%) and testing (30%) subsets. Overall, the Random Forest classifier achieves an accuracy of 97.7%. Similarly, XGboost reaches an accuracy of 97.4%. The fluidity of ongoing source association efforts offers another kind of sanity check. As our work evolved from the initial gll-psc-v23.fit 4FGL-DR2 file release to the current gll-psc-v27.fit version, five of our machine-learned subhalo candidates moved from unassociated to pulsar associations and are no longer included here.

Figure 2 compares the Receiver Operating Characteristic curves (ROC curve) for the two methods. ROC curves allow us to visualize the True Positive Rate (sensitivity) against False Positive Rate (specificity). Curves closer to the top-left corner indicate a better performance.

Having trained the classifiers using labeled LAT sources, we can now identify dark subhalo candidates among unassociated 4FGL-DR2 sources. To cut back on contamination from Galactic plane objects, we only use unassociated sources at high latitude ($|b| \geq 10^\circ$). After data cleaning, we end up with 843 unassociated sources at $|b| \geq 10^\circ$. Applying our machine-learning models to this sample, we end up with 73 4FGL-DR2 sources which both our methods predict to be dark subhalos. Table 1 presents the list of sources consistent with being subhalo candidates. Figure 3 shows their location in RA-Dec (top) and Galactic (bottom) coordinates.

Figure 1. Left: Importance of the 10 top features that were used to train the Random Forest classifier. Right: Out-of-bag performance as a function of the number of features used by the classifier.

Figure 2. ROC curve comparison of the two methods. The red thick line shows XGBoost while the black thin line shows the curve for Random Forest.

3https://fermi.gsfc.nasa.gov/ssc/data/access/lat/10yr_catalog/
44FGL J0312.1−0921, 4FGL J1221.4-0634, 4FGL J1304.4+1203, 4FGL J1400.6-1432, 4FGL J2039.4−3616
Figure 3. Locations of 73 dark subhalo candidates in the Galactic sky coordinates (empty red circles). Candidates recurring in multiple Fermi data releases are marked as solid red circles. The color-coded dots indicate the predicted distance of the GD-1 perturbing subhalo.

3 X-ray Observations

For each source in Table 1, we select all observations available in the Swift archive. The majority are part of a long-term follow-up program of unassociated Fermi sources [35]. In total, we find that 69 of the 73 sources have had their 95% error ellipse covered. For the X-ray Telescope (XRT) analysis, we used XRTDAS$^5$ within HEASOFT 6.26. We used the standard selections of $0 - 12$ in the Photon Counting (PC) mode and performed the analyses in the 0.3–10 keV energy range. In case of a point source detection within the 95% error ellipse, we extract the number of photons from a circular region around the source with radius 20 arcseconds and a 40-150 arcsecond annular region to describe the background and scale it to the source region. In total, 17 4FGL-DR2 sources show one or more X-ray detections within their 95% error ellipse. Whenever there is a detection, X-ray fluxes are calculated assuming photon index 2 and the full Galactic $N_H$ estimated using WebPIMMS$^6$. See Table 2 for flux values.

In order to place X-ray upper limits, we followed a procedure similar to [36]. We extract the number of photons from the brightest circular region within the LAT 95% ellipse. Using annular background regions, we compute the expected number of background photons and compute the 90% confidence lower limit in the background region [37]. The upper limit is derived by subtracting this lower limit from the observed photons in the source region. To determine upper limits for nondetections, we assume an absorbed power law with photon index 2 and $N_H$ from WebPIMMS. Results are listed in Table 1. In order to double check our results, we used the online Swift XRT products generator$^7$ based on tools developed for the 2SXPS Swift XRT Point Source Catalogue [38]. Finally, we also consulted the online version of the Swift-XRT Survey of Fermi Unassociated Sources program for consistency$^8$.

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5https://swift.gsfc.nasa.gov/analysis/xrt_swguide_v1_2.pdf
6https://heasarc.gsfc.nasa.gov/cgi-bin/Tools/w3pimms/w3pimms.pl
7https://www.swift.ac.uk/
8https://www.swift.psu.edu/unassociated/
4 Searching for the GD-1 Perturber

Searches for low-mass dark-matter subhalos in the Milky Way galaxy focus on detecting their effects on the motions of halo stars. Stellar streams are especially sensitive tracers of such low-amplitude gravitational perturbations [39, 40]. Formed by tidal dissolution of a progenitor star cluster or dwarf galaxy, stellar streams are long, thin and kinematically cold groups of stars moving on similar orbits through the Galaxy for billions of years [41–43]. A subhalo passing by a stellar stream imparts measurable velocity kicks to its member stars. Less massive subhalos merely increase the stream’s velocity dispersion [39], while more massive ones significantly alter stellar orbits to produce an underdensity, or a "gap", in the stellar stream [23]. Parameters of the impact, including the subhalo mass and time of impact, remain recorded in the density and kinematic structure of the perturbed region [44, 45]. Early analyses of photometric data from the Sloan Digital Sky Survey (SDSS) revealed an abundance of stream gaps consistent with perturbation by a population of CDM subhalos [46, 47]. However, deeper photometry indicates that at least some of these density variations may be due to contamination from other Milky Way stars [48], rendering the detection of dark-matter subhalos with streams inconclusive in SDSS.

Proper motions provided by the Gaia mission vastly improved the selection of stream member stars. Applied to the GD-1 stellar stream, this selection revealed significant gaps in the stream density, as well as a spur of stream stars beyond the main stream track [49]. Numerical experiments show that a recent (≈ 0.5 Gyr) impact of a massive (10^6 M_☉−10^8 M_☉) and compact (≲ 20 pc) object can produce features observed in GD-1 [26]. Even though a spur-like structure can be induced in GD-1 by the Sagittarius dwarf galaxy [27], orbital integrations in the fiducial model of the Milky Way show that none of the known objects like dwarf galaxies, globular clusters, and molecular clouds, approach GD-1 sufficiently close to reproduce the structure of the GD-1 gap-and-spur features in detail [26]. This allows the possibility that GD-1 encountered a low-mass dark-matter subhalo.

The most direct way to ascertain the nature of a stream perturber would be to observe it directly. Detailed spatial distribution (the location and size of the observed gap and spur features) and kinematics (small radial velocity offset detected between the stream and spur) of the perturbed region of GD-1 constrain the perturber’s orbit [26]. Present-day sky positions of perturbers on allowed orbits are shown in Figure 3 and color-coded by their distance from the Sun (darker colors for more nearby models). Fermi dark-matter subhalo candidates are overplotted as large red circles. The present-day positions of the GD-1 perturber were calculated assuming an axisymmetric, static, analytic model of the Milky Way’s gravitational potential, whereas the most recent dynamical studies suggest that the halo shape is spatially complex and still evolving due to ongoing mergers with Sagittarius and Large Magellanic Cloud [50, 51]. Therefore, each individual perturber position is likely associated with a systematic uncertainty stemming from an overly simplified Milky Way model. Still, positions of a number of subhalo candidates overlap the general area plausibly occupied by the GD-1 perturber and provide valuable targets for follow-up observations.

Interestingly, the distribution of Fermi subhalo candidates is asymmetric with respect to the Galactic center: most candidates are located at Galactic latitudes |b| ≲ 45°, but with a larger extension to negative Galactic longitudes, −100° ≲ l ≲ 60°. Even more strikingly, the preferred sky positions of the GD-1 perturber are along a great circle. Dynamically, their orbits appear to be aligned with the orbital plane of the Sagittarius dwarf galaxy [26]. The robustness of these asymmetries needs to be further tested, e.g., with a more comprehensive
search of the orbital parameter space in the case of the GD-1 perturber, and with a more
detailed accounting of Fermi sensitivity to subhalo signals. If confirmed, the asymmetric
distribution of subhalo candidates suggests that at least a fraction of dark matter in the
Milky Way is unrelaxed and that coherent debris flows from recent mergers [e.g., 51–54] need
to be accounted for in studies of dark matter on Galactic scales.

5 Is GD-1’s Perturbing Subhalo Detectable in Gamma Rays?

To start answering this question, we must rely on cosmological simulations and $J$
factor predictions. The expected gamma-ray flux from dark matter annihilation [55] is given by:

$$
\frac{d\phi(\Delta \Omega)}{dE_\gamma} = \left( \frac{1}{4\pi} \frac{dN_\gamma}{dE_\gamma} \langle \sigma v \rangle \right) \left( \int_{\Delta \Omega} d\Omega \int_{\text{l.o.s.}} d\ell \rho_\chi(\vec{\ell}) \right).
$$

The left term within parentheses represents the particle physics contribution where $\langle \sigma v \rangle$
is the velocity-averaged annihilation cross section, $m_{DM}$ is the WIMP mass, and $dN_\gamma/dE_\gamma$ is
the differential spectrum of gamma rays from annihilation of a pair of DM particles. The right
hand side term (as shown in parentheses) represents the $J$-factor $J$, which encloses all the
astrophysical considerations. The first integral is performed over the solid angle of a region
of interest (ROI, $\Delta \Omega$). The second integration is performed along the line of sight, and $\rho(\vec{\ell})$
is the density of the DM particles.

Finding astronomical objects that maximize $J$ is key to achieving a detection. Using
predictions from the Via Lactea II [2] dark-matter only simulation, [8] find significantly higher
$J$-factors associated to low-mass subhalos. With respect to previous works, [8] assigns “proper”
subhalo concentrations and repopulated the VL-II simulations with low-mass subhalos down
to $10^3 M_\odot$. The main limitation in these results is the unknown impact of baryons in the
subhalo population.

Possibly the most careful job assessing the gamma-ray detectability of dark matter
subhalos taking into account baryons and uncertainties in different cosmological simulations
is by [56]. Using a set of four configurations including Aquarius, Phat-ELVIS, Stref and
Lavalle, [56] find that there should be anywhere between 4 and 50 subhalos in the $\sim 10^6 – 10^8$
$M_\odot$ mass range detectable in the 3FGL catalog between 15 kpc and 160 kpc. Relative to the
3FGL catalog, the 4FGL-DR2 used in this work has twice as much exposure.

Assuming a massive $\sim 10^6 – 10^8 M_\odot$ perturbing subhalo (as in the previous section)
and a 100 GeV dark matter particle annihilating to $b\bar{b}$ at the thermal relic cross section, we
expect a 5σ significance detection of any subhalo with $J \sim 10^{17.2}$ GeV$^2$ cm$^{-5}$ within 8 kpc
[57]. If $J$ factors scale as the square of the distance to the target, one could detect subhalos
with $J \sim 10^{19.4}$ GeV$^2$ cm$^{-5}$ out to 100 kpc. From our predicted distance of the perturbing
subhalo, we find a wide distribution that would cover that detection range.

To substantiate these estimates, we use the CLUMPY code$^9$ to reproduce a scaled down
version of $J$-factor predictions here. We follow a nearly identical approach as the one used in
[58]. We generate subhalo sets ranging from $m_{\text{min}} = 10^{-6} M_\odot$ to $m_{\text{max}} = 0.01 M_{\text{tot}}$ using the
default value for CLUMPY parameters. We assume a halo mass distribution $dN/dm \propto m^{-\alpha_m}$
with $\alpha_m = 1.9$ and set a limit of 150 subhalos between $10^5 M_\odot$ and $10^{10} M_\odot$. An Einasto
profile is the default configuration for the subhalo profile. In Figure 4, we show the brightest
2500 subhalos in 5 realizations at high Galactic latitude ($|b| \geq 10^\circ$). As expected, we find

$^9$https://clumpy.gitlab.io/CLUMPY/
some of the largest J-factors in the $\sim 10^6 - 10^8 M_\odot$ subhalo mass range within distances between 1 and 100 kpc. For a complete analysis using 500 simulations and a more complete discussion of CLUMPY caveats see [58].

6 They Might Be Pulsars

Throughout our discussion, we have noted that it might not be possible to distinguish between subhalos from annihilating $\sim 20$–$70$ GeV dark matter particles and gamma-ray pulsars using spectral information alone. Based on population syntheses of millisecond pulsars, some models predict that Fermi could detect around 170 MSPs within 10 years [59]. As of this writing, there are 118 MSPs in the Public List of LAT-Detected Gamma-Ray Pulsars 10. The upcoming Third Fermi LAT Pulsar Catalog (3PC in preparation) [60] will significantly expand the census of gamma-ray pulsars and facilitate a proper comparison with the list of candidates presented here.

7 Conclusions and Future Work

We present a list of 73 4FGL-DR2 sources that appear to be promising dark matter subhalo candidates. Out of the 69 covered by Swift, 17 show at least one X-ray source within the LAT ellipse. If we only consider gamma-ray sources detected in more than one catalog (1FGL, 2FGL, 3FGL, 4FGL, 4FGL-DR2), one can narrow the list further to 40 sources. From the analysis presented here and recurrent selection as subhalo candidates in previous works [8, 18, 61], we highlight a subset of seven sources that should be further investigated as potential subhalo candidates: 4FGL J0545.7+6016, 4FGL J0953.6-1509, 4FGL J1225.9+2951, 4FGL J1539.4-3323, 4FGL J1543.6-0244, 4FGL J2112.5-3043, and 4FGL J2133.1-6432. Out

10https://confluence.slac.stanford.edu/display/GLAMCOG/Public+List+of+LAT-Detected+Gamma-Ray+Pulsars
of these seven candidates, 4FGL J1539.4-3323 and 4FGL J2112.5-3043 lie in sky locations with the highest density of allowed GD-1 perturber models with predicted distances smaller than 40 kpc.

It is important to note that the presumed subhalo spectra from annihilating $\sim 20–70$ GeV dark matter particles is consistent with the intriguing Galactic Center signal [62, 63]. Now there are two general caveats when interpreting our results. First, any correspondence between simulations and 4FGL-DR2 subhalo locations as in Figure 3 should not be taken directly as a distance proxy. Second, the mere presence of X-ray point sources within a LAT 95% error ellipse does not necessarily translate into an actual source association [28]. However, X-ray sources can provide initial follow-up targets for optical spectroscopy.

As the characterization of stellar streams improves, it might be possible to actually pinpoint the exact location of dark matter perturbing subhalos directly. This will further narrow down the potential gamma-ray counterpart. This work is a first attempt at what might be possible to achieve in the future.

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| Source name | RA        | Dec       | P(RF) | P(XGboost) | Initial Detection |
|-------------|-----------|-----------|-------|------------|------------------|
| 4FGL J0036.6+3059 | 0.9045    | 30.9898   | 0.88  | 0.94       | 3FGL             |
| 4FGL J0048.6−6347 | 12.1685   | -63.7914  | 0.88  | 0.91       | 2FGL             |
| 4FGL J0139.5−2228 | 24.8971   | -22.4777  | 0.54  | 0.66       | 4FGL             |
| 4FGL J0336.0+7502 | 54.0249   | 75.05     | 0.99  | 0.99       | 1FGL             |
| 4FGL J0341.9+3153c | 55.4852   | 31.8952   | 0.85  | 0.99       | 2FGL             |
| 4FGL J0414.7−4300 | 63.6977   | -43.012   | 0.90  | 0.96       | 4FGL             |
| 4FGL J0418.9+6636 | 64.7254   | 66.6001   | 1.00  | 1.00       | 2FGL             |
| 4FGL J0436.9−2915 | 69.2464   | 29.2565   | 0.72  | 0.63       | 4FGL             |
| 4FGL J0533.6+5945 | 83.4175   | 59.7622   | 0.98  | 0.97       | 3FGL             |
| 4FGL J0545.7+6016 | 86.4419   | 60.2704   | 0.95  | 0.99       | 1FGL             |
| 4FGL J0802.1−5612 | 120.5456  | -56.2012  | 1.00  | 0.99       | 1FGL             |
| 4FGL J0906.8−2122 | 136.7046  | -21.3724  | 0.66  | 0.84       | 4FGL             |
| 4FGL J0940.3−7610 | 145.0989  | -76.1794  | 0.99  | 0.96       | 1FGL             |
| 4FGL J1106.7−1742 | 166.699   | -17.7148  | 0.83  | 0.68       | 1FGL             |
| 4FGL J1120.0−2204 | 170.0016  | -22.0779  | 0.99  | 0.96       | 1FGL             |
| 4FGL J1126.0−5007 | 171.5145  | -50.1194  | 0.74  | 0.69       | 3FGL             |
| 4FGL J1204.5−5032 | 181.1483  | -50.5456  | 0.98  | 0.91       | 4FGL             |
| 4FGL J1207.4−4536 | 181.8734  | -45.6125  | 0.99  | 0.88       | 3FGL             |
| 4FGL J1225.9−2951 | 186.4902  | -29.8596  | 1.00  | 0.97       | 1FGL             |
| 4FGL J1331.6−5116 | 187.9102  | -51.2672  | 0.99  | 0.99       | 1FGL             |
| 4FGL J1345.9−2612 | 206.4815  | -26.2116  | 0.87  | 0.63       | 2FGL             |
| 4FGL J1400.0−2415 | 210.0206  | -24.266   | 0.94  | 0.99       | 2FGL             |
| 4FGL J1429.8−0739 | 217.4512  | -7.6506   | 0.52  | 0.69       | 4FGL             |
| 4FGL J1458.8−2120 | 224.7033  | -21.3388  | 0.85  | 0.99       | 2FGL             |
| 4FGL J1526.6−2743 | 231.6709  | -27.7327  | 1.00  | 0.98       | 4FGL             |
| 4FGL J1526.6−3810 | 231.6594  | -38.169   | 0.99  | 0.99       | 3FGL             |
| 4FGL J1527.8+1013 | 233.3757  | 10.2286   | 0.91  | 0.93       | 4FGL             |
| 4FGL J1530.0−1522 | 232.5011  | -15.3759  | 0.96  | 0.98       | 4FGL             |
| 4FGL J1539.4−3323 | 234.8511  | -33.3987  | 0.85  | 0.99       | 1FGL             |
| 4FGL J1543.6−0244 | 235.9077  | -2.7471   | 0.96  | 0.56       | 1FGL             |
| 4FGL J1544.2−2554 | 236.0523  | -25.9125  | 1.00  | 0.98       | 2FGL             |
| 4FGL J1602.2+2305 | 240.5576  | 23.0969   | 0.93  | 0.99       | 2FGL             |
| 4FGL J1612.1+1407 | 243.0313  | 14.1168   | 0.90  | 0.99       | 2FGL             |
| 4FGL J1622.2−7202 | 245.5525  | -72.0399  | 0.89  | 0.99       | 4FGL             |
| 4FGL J1623.9−6036 | 245.9865  | -69.6069  | 0.81  | 0.87       | 4FGL             |
| 4FGL J1630.1−1049 | 247.5289  | -10.8183  | 0.99  | 0.74       | 1FGL             |
| 4FGL J1646.7−2154 | 251.6851  | -21.9075  | 0.97  | 0.98       | 4FGL             |
| 4FGL J1656.4−0410 | 254.1194  | -4.1702   | 0.78  | 0.99       | 4FGL             |
| 4FGL J1659.0−0140 | 254.7655  | -1.6775   | 0.75  | 0.99       | 2FGL             |
| 4FGL J1700.0−0122 | 255.0191  | -1.3734   | 0.82  | 0.70       | 4FGL             |
| 4FGL J1709.9−0900 | 257.4834  | -9.0144   | 0.96  | 0.99       | 4FGL             |
| 4FGL J1711.9−1922 | 257.9776  | -19.3676  | 0.96  | 0.99       | 4FGL             |
| Name                | RA        | Dec       | FWHM      | LL         | Type |
|---------------------|-----------|-----------|-----------|------------|------|
| 4FGL J1717.5−5804   | 269.3784  | −58.0706  | 0.98      | 0.92       | 2FGL |
| 4FGL J1720.6+0708   | 260.1638  | 7.1469    | 0.99      | 0.65       | 1FGL |
| 4FGL J1722.8−0418   | 260.7115  | −4.3033   | 0.72      | 0.74       | 1FGL |
| 4FGL J1730.4−0359   | 262.6086  | −3.9923   | 1.00      | 0.99       | 1FGL |

Table 1: Subhalo candidates in the 4FGL-DR2
| Source name          | $N_H$ (cm$^2$) | X-ray Flux Upper Limit (erg cm$^{-2}$ s$^{-1}$) |
|----------------------|----------------|-----------------------------------------------|
| 4FGL J003.6+3059     | $4.9 \times 10^{20}$ | $1.5 \times 10^{-14}$ (PS)                    |
| 4FGL J004.6−6347     | $2.2 \times 10^{20}$ | $< 8.4 \times 10^{-14}$                       |
| 4FGL J013.9−2228     | $9.9 \times 10^{19}$ | $< 8.9 \times 10^{-14}$                       |
| 4FGL J033.6+7502     | $1.5 \times 10^{21}$ | $4.5 \times 10^{-14}$ (PS)                    |
| 4FGL J034.9+3153c    | $1.2 \times 10^{21}$ | $8.6 \times 10^{-13}$ (PS)                    |
| 4FGL J043.9+2915     | $1.7 \times 10^{21}$ | $< 1.6 \times 10^{-13}$                      |
| 4FGL J044.7+2446     | $2.1 \times 10^{21}$ | $< 8.6 \times 10^{-13}$                      |
| 4FGL J053.6+5945     | $1.8 \times 10^{21}$ | $< 1.2 \times 10^{-13}$                      |
| 4FGL J054.7+6016     | $1.5 \times 10^{21}$ | $< 2.0 \times 10^{-13}$                      |
| 4FGL J080.2−5612     | $1.4 \times 10^{21}$ | $4.1 \times 10^{-14}$ (PS)                    |
| 4FGL J090.6−2122     | $1.1 \times 10^{21}$ | $< 8.1 \times 10^{-14}$                      |
| 4FGL J094.0−7610     | $9.6 \times 10^{20}$ | $2.2 \times 10^{-14}$ (PS)                    |
| 4FGL J095.6−1509     | $5.8 \times 10^{19}$ | $< 7.7 \times 10^{-14}$                      |
| 4FGL J110.6−1742     | $4.0 \times 10^{20}$ | Partially Covered                             |
| 4FGL J112.0−2204     | $4.1 \times 10^{20}$ | $7.2 \times 10^{-14}$ (PS)                    |
| 4FGL J112.6−5007     | $9.1 \times 10^{20}$ | $6.2 \times 10^{-14}$ (PS) (also GRB 140719A in error ellipse) |
| 4FGL J120.5−5032     | $1.1 \times 10^{21}$ | $< 7.5 \times 10^{-14}$                      |
| 4FGL J120.7−4536     | $6.7 \times 10^{20}$ | $2.8 \times 10^{-14}$ (PS)                    |
| 4FGL J122.9−2951     | $1.6 \times 10^{20}$ | $< 8.7 \times 10^{-14}$                      |
| 4FGL J123.6−5116     | $1.4 \times 10^{21}$ | $< 1.2 \times 10^{-13}$                      |
| 4FGL J134.5−2612     | $4.7 \times 10^{20}$ | $< 8.6 \times 10^{-14}$                      |
| 4FGL J140.0−2415     | $5.1 \times 10^{20}$ | $< 6.4 \times 10^{-14}$                      |
| 4FGL J142.9−0739     | $4.6 \times 10^{20}$ | $8.5 \times 10^{-14}$ (PS)                    |
| 4FGL J145.8−2120     | $7.4 \times 10^{20}$ | $< 8.3 \times 10^{-14}$                      |
| 4FGL J152.6−3810     | $1.0 \times 10^{21}$ | $< 1.4 \times 10^{-13}$                      |
| 4FGL J152.6−2743     | $7.7 \times 10^{20}$ | $< 1.4 \times 10^{-13}$                      |
| 4FGL J152.8+1013     | $2.8 \times 10^{20}$ | $1.1 \times 10^{-13}$ (PS)                    |
| 4FGL J153.0−1522     | $8.3 \times 10^{20}$ | $1.4 \times 10^{-13}$ (PS)                    |
| 4FGL J153.9−3323     | $9.0 \times 10^{20}$ | $< 9.0 \times 10^{-13}$                      |
| 4FGL J154.3−0244     | $7.9 \times 10^{20}$ | $< 9.6 \times 10^{-14}$                      |
| 4FGL J154.4−2554     | $1.1 \times 10^{21}$ | $< 9.9 \times 10^{-14}$                      |
| 4FGL J160.2+2305     | $4.3 \times 10^{20}$ | $< 1.2 \times 10^{-13}$                      |
| 4FGL J161.2+1407     | $3.1 \times 10^{20}$ | $< 8.4 \times 10^{-14}$                      |
| 4FGL J162.2−7202     | $7.5 \times 10^{20}$ | $8.9 \times 10^{-14}$ (PS)                    |
| 4FGL J162.3−6936     | $9.7 \times 10^{20}$ | $< 7.7 \times 10^{-14}$                      |
| 4FGL J162.7−5749     | $2.7 \times 10^{21}$ | Partially Covered                             |
| 4FGL J163.0−1049     | $1.3 \times 10^{21}$ | $< 1.1 \times 10^{-13}$                      |
| 4FGL J164.6−2154     | $1.3 \times 10^{21}$ | Not Covered                                   |
| 4FGL J165.4−0410     | $1.0 \times 10^{21}$ | $1.2 \times 10^{-14}$ (PS)                    |
| 4FGL J165.9−0140     | $8.4 \times 10^{20}$ | $< 1.3 \times 10^{-13}$                      |
| 4FGL J170.0−0122     | $8.5 \times 10^{20}$ | $2.6 \times 10^{-13}$ (PS)                    |
| 4FGL J170.9−0900     | $1.5 \times 10^{21}$ | $< 1.1 \times 10^{-13}$                      |
| 4FGL J171.9−1922     | $1.5 \times 10^{21}$ | $< 1.9 \times 10^{-13}$                      |
| 4FGL J171.7−5804     | $9.2 \times 10^{20}$ | $< 2.6 \times 10^{-13}$                      |

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| Source                  | X-ray Limit   | Swift Coverage   |
|------------------------|---------------|------------------|
| 4FGL J1720.6+0708      | $6.2 \times 10^{20}$ | $< 1.0 \times 10^{-13}$ |
| 4FGL J1722.8-0418      | $1.3 \times 10^{21}$ | $< 1.0 \times 10^{-13}$ |
| 4FGL J1730.4-0359      | $1.5 \times 10^{21}$ | $< 9.0 \times 10^{-13}$ |
| 4FGL J1757.7-6032      | $6.8 \times 10^{20}$ | $< 1.2 \times 10^{-13}$ |
| 4FGL J1813.5+2819      | $7.4 \times 10^{20}$ | $< 1.2 \times 10^{-13}$ |
| 4FGL J1818.6+1316      | $9.8 \times 10^{20}$ | $< 1.2 \times 10^{-13}$ |
| 4FGL J1823.2+1209      | $1.3 \times 10^{21}$ | $< 1.3 \times 10^{-13}$ |
| 4FGL J1824.2-5427      | $7.7 \times 10^{20}$ | $< 1.3 \times 10^{-13}$ |
| 4FGL J1827.5+1141      | $1.4 \times 10^{21}$ | $< 1.1 \times 10^{-13}$ |
| 4FGL J1831.1-6503      | $6.1 \times 10^{20}$ | $< 1.1 \times 10^{-13}$ |
| 4FGL J1842.1+2737      | $1.1 \times 10^{21}$ | $< 1.1 \times 10^{-13}$ |
| 4FGL J1845.8-2521      | $1.2 \times 10^{21}$ | $< 1.3 \times 10^{-13}$ |
| 4FGL J1855.6-3603      | $6.4 \times 10^{20}$ | $< 1.4 \times 10^{-13}$ |
| 4FGL J1858.3-5424      | $6.0 \times 10^{20}$ | $< 1.8 \times 10^{-13}$ |
| 4FGL J1906.4-1757      | $9.1 \times 10^{20}$ | Not covered      |
| 4FGL J1910.7-5320      | $5.3 \times 10^{20}$ | $< 3.3 \times 10^{-13}$ |
| 4FGL J1920.0-2622      | $7.9 \times 10^{20}$ | $< 1.2 \times 10^{-13}$ |
| 4FGL J1924.8-1035      | $1.2 \times 10^{21}$ | $< 1.1 \times 10^{-13}$ |
| 4FGL J1949.2-1453      | $7.8 \times 10^{20}$ | $< 1.1 \times 10^{-13}$ |
| 4FGL J2026.3+1431      | $6.2 \times 10^{20}$ | $< 1.0 \times 10^{-13}$ |
| 4FGL J2029.5-4237      | $3.7 \times 10^{20}$ | $3.7 \times 10^{-13}$ (PS) |
| 4FGL J2043.9-4802      | $2.7 \times 10^{20}$ | $< 5.2 \times 10^{-14}$ |
| 4FGL J2112.5-3043      | $6.4 \times 10^{20}$ | $< 2.5 \times 10^{-14}$ |
| 4FGL J2133.1-0432      | $3.2 \times 10^{20}$ | $< 6.0 \times 10^{-14}$ |
| 4FGL J2212.4+0708      | $6.2 \times 10^{20}$ | $1.3 \times 10^{-13}$ (PS) |
| 4FGL J2219.7-6837      | $2.8 \times 10^{20}$ | $< 1.5 \times 10^{-13}$ |
| 4FGL J2250.5+3305      | $8.0 \times 10^{20}$ | $2.0 \times 10^{-13}$ (PS) |

**Table 2:** X-ray upper limits for 4FGL-DR2 subhalo candidates. PS denotes at least one X-ray point source within the LAT error ellipse. Partial or no **Swift** coverage are also indicated.