EXTREME ULTRAVIOLET SPECTRA OF SOLAR FLARES FROM THE EXTREME ULTRAVIOLET SPECTROHELIOGRAPH SPIRIT ONBOARD THE CORONAS-F SATELLITE

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ABSTRACT
We present detailed extreme ultraviolet (EUV) spectra of four large solar flares: M5.6, X1.3, X3.4, and X17 classes in the spectral ranges 176–207 Å and 280–330 Å. These spectra were obtained by the slitless spectroheliograph SPIRIT onboard the CORONAS-F satellite. To our knowledge, these are the first detailed EUV spectra of large flares obtained with a spectral resolution of ~0.1 Å. We performed a comprehensive analysis of the obtained spectra and provide identification of the observed spectral lines. The identification was performed based on the calculation of synthetic spectra (the CHIANTI database was used), with simultaneous calculations of the differential emission measure (DEM) and density of the emitting plasma. More than 50 intense lines are present in the spectra that correspond to a temperature range of $T = 0.5–16$ MK, most of the lines belong to Fe, Ni, Ca, Mg, and Si ions. In all the considered flares, intense hot lines from Ca xii, Ca xiii, Fe xx, Fe xxii, and Fe xxiv are observed. The calculated DEMs have a peak at $T \sim 10$ MK. The densities were determined using Fe xi–Fe xiii lines and averaged $6.5 \times 10^9$ cm$^{-3}$. We also discuss the identification, accuracy, and major discrepancies of the spectral line intensity prediction.

Key words: Sun: activity – Sun: corona – Sun: flares – Sun: UV radiation

Online-only material: color figures

1. INTRODUCTION
The extreme ultraviolet (EUV) emission of the solar corona has been studied since the beginning of the space era due to the rich informational content of the registered spectra. Analysis of such spectra allows the determination of various plasma characteristics, such as temperature and density, and provides information about dynamic processes that take place in the solar corona. In addition, the EUV spectra of different coronal phenomena have become a subject of interest in a number of different areas such as atomic physics, astrophysics, and physics of plasmas.

Numerous spectroscopic observations have been carried out using spectroscopic instruments of different types: slit spectrographs with high spatial resolution, such as SERTS (Neupert et al. 1992), Solar and Heliospheric Observatory/Coronal Diagnostic Spectrometer (SOHO/CDS; Harrison et al. 1995), Hinode/Extreme-Ultraviolet Imaging Spectrometer (EIS; Culhane et al. 2007), spectroheliographs with imaging capabilities such as S082A/Skylab (Tousey et al. 1977), SPIRIT/CORONAS-F (Zhitnik et al. 2002), and full-Sun spectrographs, which obtain spectra from the whole solar disk, such as those on the Aerobee rocket (Malinovsky & Heroux 1973) or the Solar Dynamics Observatory/Extreme Ultraviolet Variability Experiment (SDO/EVE; Woods et al. 2012).

Data obtained in these experiments have been used for various goals such as for development of atlases of spectral lines, validation of atomic data, measurement of temperature and density of the emitting plasma in different structures, and the determination of presence of upflows or downflows, etc. Among the structures that were studied, there are quiet Sun (QS) regions (Brosius et al. 1996), active regions (AR), cores (Tripathi et al. 2011), off-limb AR plasma (O’Dwyer et al. 2011), AR mosses (Tripathi et al. 2010), coronal streamers (Parenti et al. 2003), bright points (Ugarte-Urra et al. 2005), and others.

Whereas solar flares have also been observed by spectrographs, obtaining EUV spectra of solar flares is not so common. The first systematic analysis of EUV flaring spectra was presented by Dere (1978). The author analyzed more than 50 photographic plates from the S082A spectroheliograph on Skylab and constructed a catalog of spectral lines in the range 171–630 Å. The catalog included relative intensities of more than 200 spectral lines.

Systematic studies of EUV spectra of solar flares have been continued on subsequent satellites: SOHO (launched in 1995), Hinode (launched in 2006), and SDO (launched in 2010). The CDS spectrograph onboard the SOHO satellite registered several large solar flares during their decay phases. The first analysis of a CDS flare was made by Czaykowska et al. (1999). The authors analyzed intensities of spectral lines during the decay phase of an M6.8 flare and determined the density and temperature of the post-flare loops. Del Zanna et al. (2006) also performed an analysis of spectra of a X17 flare during the decay phase. The authors studied Doppler shifts and found them to be consistent with those predicted by a simple hydrodynamics model. It is worth noting that due to the telemetry constraints of CDS, all these flares were observed in a fast-rastering regime in only six narrow spectral windows, covering only a small portion of the wide spectral ranges 308–381 and 513–633 Å of the CDS.

The EIS spectrograph onboard the Hinode satellite used an improved optical layout with high-efficiency EUV optics and detectors. Therefore, EIS has superb spectral, spatial, and temporal resolution, as well as higher telemetry volumes, which allow spectra to be investigated in much higher detail, such as with a wider set of spectral lines, higher cadence, and higher spatial resolution. EIS has observed a large span of flares, starting from small B2 class flares (Del Zanna et al. 2011) to large M1.8 class flares (Doschek et al. 2013). However,
despite all its advantages, EIS usually observes flares in a coarse rastering regime. This fact limits the number of spectral lines observed in a flare; for example, Watanabe et al. (2010) used only 17 lines for plasma diagnostics from the whole spectral range 170–210 and 250–290 Å. There is a case when EIS registered a full CCD flare spectrum (Doschek et al. 2013); however, the authors focused on Doppler shift analysis and used only 17 out of the 500 lines registered by EIS.

The EVE spectrometer onboard the SDO builds whole-Sun spectra in the range 10–1050 Å. It has a moderate spectral resolution of 1 Å but operates with an unprecedented 10 s cadence and almost a 100 % duty cycle. There are two main difficulties in the analysis of the EVE spectra: it has no spatial resolution—the flare spectrum is mixed with the spectrum from the rest of the Sun and, due to moderate spectral resolution of EVE, most of the lines are blended. Despite these obstacles, EVE is widely used in solar investigations: for studies of the thermal evolution of flaring plasma (Chamberlin et al. 2012), Doppler shift studies (Hudson et al. 2011), and high temperature plasma electron density diagnostics (Milligan et al. 2012).

Without diminishing the importance of the information obtained in these experiments, it should be noted that a small number of EUV spectra of solar flares have been registered so far and published catalogs of spectral lines are limited.

In this paper, we take advantage of the SPIRIT EUV spectroheliograph and perform a comprehensive analysis of EUV spectra of four large solar flares. The flares of M5.6, X1.3, X3.4, and X17 classes have been observed by the slitless EUV spectroheliograph SPIRIT onboard the CORONAS-F satellite. The spectroheliograph operated in two wavelength ranges of 176–207 and 280–330 Å and had a spectral resolution of 0.1 Å. We perform an absolute calibration of SPIRIT spectral fluxes using simultaneous SOHO/Extreme Ultraviolet Imaging Telescope (EIT) images. In order to identify the obtained spectra, we use an original approach, based on calculation of synthetic spectra and their subsequent modification to match the observational data. Simultaneously, we calculate the differential emission measure (DEM) and $n_e$ of the emitting plasma and repeat iteratively the whole procedure of identification several times.

We provide identification of more than 50 spectral lines in each spectral band for each flare. In addition to spectral line intensities, we calculate the DEM and plasma density for each flare. The obtained information can be used not only for modeling spectral fluxes in different EUV spectral bands and for refinement of the atomic data, but also for studying the flares themselves and validating models for flare plasma evolution.

The obtained spectra, synthetic spectra, DEMs, and proposed IDL software are available at http://coronas.izmiran.ru/F/SPRINT/ or by request from S. Shestov.

2. OBSERVATIONS

The SPIRIT complex of instrumentation was launched onboard the CORONAS-F satellite (Oraevsky & Sobelman 2002) on 2001 July 31 from Plesetsk cosmodrome in northern Russia. The satellite was placed on a near-polar orbit with an inclination of 82° and a perigee of 500 km. The satellite carried 12 scientific instruments for the measurement of both particle and electromagnetic emission of the Sun. The SPIRIT instrumentation was developed in the Lebedev Physical Institute of the Russian Academy of Sciences and consisted of telescopic and spectroheliographic channels for observation of the solar corona in different soft X-ray and EUV spectral bands (Zhitnik et al. 2002).

The EUV spectroheliograph SPIRIT consisted of two similar independent spectral channels: the V190 channel for the 176–207 Å range and the U304 channel for the 280–330 Å range. Both channels were built using a slitless optical scheme (see Figure 1). The solar EUV emission enters through an entrance filter and falls on a diffraction grating (with a grazing angle $\varphi \sim 1:5$). The diffracted radiation is focused on a detector by a mirror with a multilayer coating.

The slitless optical scheme observes the full-Sun field of view on the detector, which allowed us to obtain as many as 30 spectroheliograms with large solar flares over 4.5 yr of the satellite’s lifetime.

For the analysis, we have selected the following flares: an M5.6 observed on 2001 September 16, an X3.4 observed on 2001 December 28, an X1.3 observed on 2004 July 16, and an X17 observed on 2005 September 7. All these flares are long-duration events (LDEs), cover a broad range of flare intensity, and have been registered in different phases of their decay. The
X-ray light curves of the flares measured by GOES are shown in Figure 2. Each SPIRIT spectroheliogram was obtained in a single exposure (the exposures are denoted by vertical lines in Figure 2). The exposure times for the M5.6 and X3.4 flares were 37 s; the exposure times were 150 s for the X1.3 and X17 flares. Some details of the analyzed flares are given in Table 1.

### Table 1

| Flares Class | Date       | GOES Peak Time (UT) | SPIRIT Obs. Start (UT) | NOAA AR | Type, $t_{\text{decay}}$ |
|--------------|------------|---------------------|------------------------|---------|-------------------------|
| M5.6         | 2001 Sep 16 | ~03:50              | 03:59:36               | 9608    | LDE, 1 hr 30 minutes    |
| X1.3         | 2004 Jul 16 | ~02:05              | 02:07:54               | 10649   | LDE, ~7 hr              |
| X3.4         | 2001 Dec 28 | ~20:40              | 21:21:45               | 9767    | LDE, 2 hr 40 minutes    |
| X17          | 2005 Sep 7  | ~17:40              | 20:04:22               | 10808   | LDE, 5 hr 50 minutes    |

Comparison of the spectroheliograms and the extracted spectra shows that emission of “cold” coronal lines (like Si IX, Mg VIII, Fe XI, and Fe XII with $T_{\text{max}} \sim 1$–2 MK) originates from the whole solar disk, but due to contraction these monochromatic full-disk images look like ellipses. Emission of “hot” coronal lines (Ca XVII, Ca XVIII, Fe XX, Fe XXII, and Fe XXIV with $T_{\text{max}} > 6$ MK) is produced mainly in flaring regions, which correspond to bright points in the spectroheliograms.

The interpretation of the spectroheliograms involves the following steps: (1) obtaining spectra of a particular region and determining the wavelength scale, (2) subtracting the background from the spectra, and (3) identifying spectral lines with a subsequent analysis of the spectral data.

For obtaining spectra from the spectroheliograms, we have developed IDL software that implements a geometrical model of the spectroheliograph. According to the model, for a particular point source, the position on the CCD detector is calculated using its solar coordinates, wavelengths, and several parameters (such as direction to the solar center, groove density of the diffraction grating, focal length and direction of the focusing mirror, relative position of the CCD detector, etc.). The geometrical model automatically takes into account contraction of solar disk images and nonlinear wavelength scale across the CCD detector. Thus, to obtain spectra of a particular region and calculate the wavelength scale for it, one has only to point to the region on the solar disk. The accuracy of the obtained wavelength scale is comparable to the spectral size of 1 pixel ($\sim 0.04$ Å).
For background subtraction, we used a procedure similar to that of Thomas & Neupert (1994)—we interpolated values outside spectral lines and subtracted the interpolation from spectra.

Before the identification, we also carefully removed the strong Si xi ($\lambda = 303.33$ Å) and He ii (doublet $\lambda = 303.78 + 0.79$ Å) blend from spectra. This reveals the spectral lines of Ca xviii, Ni xiv, Fe xv ($\lambda \approx 302$ Å) and Fe xvii, and Fe xv ($\lambda \approx 305$ Å), which lie on the wings of the Si xi/He ii blend. These lines are well distinguished on the wings of the blend (see Figure 3); therefore, we remove the wings of the blend by interpolating the values outside the lines and manually zero out the core of the blend.

In order to identify the observational spectrum and measure intensities of separate spectral lines, we produced a synthetic spectrum that fits the observational data. To produce a synthetic spectrum, we use transitions and wavelengths from CHIANTI (CHIANTI v.6; Dere et al. 1997, 2009), set the line widths $\sigma$ in accordance with the instrumental FWHM ($\sigma = -0.201 + 1.43 \times 10^{-3} \times \lambda$ (Å) for the V190 channel and $\sigma = 0.1$ (Å) for the U304 channel), and vary the intensities to match the observational data. However, straightforward fitting is not possible due to the relatively low spectral resolution of SPIRIT—$\sigma \sim 0.1$ Å and blending of most of the lines.

We overcome this obstacle using an iterative procedure (see Figure 5), which consists of an initial step: measurement of intensities of a small number of spectral lines and calculation of plasma parameters—DEM and $n_e$ (see Shestov et al. 2009, 2010) and further (iterative) steps.

1. Calculation of the synthetic spectra.
2. Automated adjustment of spectral line intensities to match the observational data. During the adjustment, the ratio of the blended lines is kept constant.
3. Manual adjustment of intensities of particular spectral lines. Using the DEM and \( n_e \) analysis data, we adjust intensities of blended lines to reach a better agreement with theory (reducing \( \chi^2 \) in the DEM reconstruction and compliance with other lines in the L-function analysis—see below).

4. Calculation of DEM and \( n_e \).

The larger number of spectral lines used for analysis during iterative steps almost completely eliminates errors due to possible misidentification or other errors. The iterative procedure turned out to be fast and stable—after the second step there are no considerable changes in DEMs and synthetic spectra. So, in our approach, the plasma diagnostic was an essential part of the line identification—we used plasma parameters to resolve blended lines.

For the calculation of synthetic spectra, we used the standard CHIANTI procedures ch_synthetic and make_chianti_spec, coronal abundances sun_coronal.abund, and chianti.ioneq ionization equilibrium.

For the DEM reconstruction, we used a genetic algorithm (GA; Siarkowski et al. 2008). The algorithm is based on ideas of biological evolution and natural selection. It starts from randomly chosen initial populations of different DEMs and produces a new generation of DEMs by crossover and mutations. The procedure stops when a local \( \chi^2 \) minimum is found. The peculiar feature of the method is that since it is based on a random evolution, different runs of the procedure on a single data set give different (but similar) results. The discrepancy among different runs directly shows the confidence of the DEM reconstruction.
For the DEM analysis, we carefully chose 46 spectral lines (Table 2)—almost all strong spectral lines, observed by SPIRIT. The exceptions are Fe xvi 284.16 Å and the Si xi/He ii blend with λ ∼ 304 Å. Both of these lines are very intense, which is likely to cause saturation of the SPIRIT detector. Also, the observed intensity of the Fe xvi line shows a systematic discrepancy with other Fe xvi lines (we will discuss possible reasons later). Nevertheless, the spectral lines analyzed cover a wide temperature range—from $T_{\text{max}} \sim 1$ MK (Mg viii) to $T_{\text{max}} \sim 16$ MK (Fe xxiv). Using a large number of lines almost completely eliminates the sensitivity of the reconstructed DEM to the intensity of a particular line, improving the reliability of the reconstruction.

Electron density $n_e$ was obtained using a modified L-function analysis (Landi & Landini 1998). According to the original method proposed by the authors, L-functions of all spectral lines of a particular ion should intersect at a single point, corresponding to the density of the emitting plasma. The L-function of a spectral line is defined as a ratio of the measured intensity over the contribution function $G(T_e, n_e)$, plotted as a function of density. We slightly simplify the definition of the L-function by using $T_{\text{max}}$ instead of $T_e$ (a specially computed temperature) and plot the L-functions for major lines of the Fe xi, Fe xii, Fe xiii, Fe xv, Mg viii, and Ni xvi ions.

3.2. Absolute Calibration of SPIRIT Fluxes

No absolute ground calibration was carried out before the launch of SPIRIT. A lack of calibration cripples the spectroheliographic diagnostic capabilities. However, the spectral ranges of the SPIRIT V190 and U304 channels overlap with spectral responses of the EIT 195 Å and 304 Å channels and it is possible to cross-calibrate the SPIRIT data with the EIT data.

The total flux $F$ in an EIT image expressed in digital numbers (dn) can be expressed as

$$F = \int s(\lambda) b(\lambda) d\lambda, \quad (1)$$

where $s(\lambda)$ is real incident spectral flux in units of erg s$^{-1}$ cm$^{-2}$ Å$^{-1}$ and $b(\lambda)$ is the EIT spectral sensitivity, expressed in units of cm$^2$ dn erg$^{-1}$ and obtained with the eit_parm function from Solar Software. $s(\lambda)$ can be expressed as

$$s(\lambda) = k \cdot i(\lambda), \quad (2)$$

where $i(\lambda)$ in dn is spectral flux measured by SPIRIT and $k$ erg s$^{-1}$ cm$^{-2}$ Å$^{-1}$ dn$^{-1}$ is the calibration coefficient to be found. From Equations (1) and (2), we can calculate $k$:

$$k = \frac{F}{\int i(\lambda)b(\lambda)d\lambda}. \quad (3)$$

The relative spectral flux $i(\lambda)$ was obtained by integrating the whole SPIRIT spectroheliogram along the spatial axis. The total EIT flux $F$ was obtained by integrating the whole EIT image (195 Å for the V190 channel and 304 Å for the U304 channel). The V190 channel spectroheliogram containing the M5.6 flare and the whole-Sun relative spectral flux $i(\lambda)$, both multiplied by $b(\lambda)$, is given in Figure 6.

We carried out this procedure for all flare spectra presented in this work and converted the spectra into physical units. However, we believe that the calibration coefficient $k$ obtained for the U304 channel is less reliable than that for the V190 channel due to the possible nonlinear response of the SPIRIT detector to the intense fluxes. That is why we performed an independent verification of the obtained absolute fluxes. The verification uses a spectroscopic approach and consists of the following: during the DEM calculation, the $\chi^2$ parameter is minimized. We introduced a calibration correction factor $\kappa$ for the U304 channel and calculated $\chi^2$ values for a range of $\kappa$ values. The minimum $\chi^2$ value gives best cross-calibration $\kappa$ from the spectroscopic point of view. The calculated best $\kappa$ values are 1.0, 0.63, and 1.1 for the M5.6, X1.3, and X17 flares, respectively. These $\kappa$ have been taken into account—we modified the data in the U304 channel spectra.

4. RESULTS

We have analyzed spectra of the four flares and note three main results of our analysis:

1. a catalog of EUV spectral lines observed in large solar flares;
2. DEM and $n_e$ of the emitting plasma; and
3. a benchmark of the atomic database, by analyzing ratios of the observed and calculated spectral line intensities.

4.1. Catalog of Spectral Lines

Comparisons of observational and fitted spectra are given in Figure 7 (V190) and Figure 8 (U304). The black curve denotes
Figure 6. Top panel: the V190 spectroheliogram registered on 2001 September 9, convolved with the EIT 195 Å bandpass function. Bottom panel: solid line—SPIRIT spectra from the aforementioned spectroheliogram integrated along the spatial axis and multiplied by the EIT 195 Å bandpass function. Dashed line—normalized EIT 195 Å bandpass function.

(A color version of this figure is available in the online journal.)

Table 2

| N | Ion | λ (Å) | log Tm (K) | N | Ion | λ (Å) | log Tm (K) | N | Ion | λ (Å) | log Tm (K) |
|---|-----|-------|------------|---|-----|-------|------------|---|-----|-------|------------|
| 1 | Fe xi | 180.41 | 6.2 | 17 | Fe xiii | 196.54 | 6.3 | 33 | Fe xiii | 312.11 | 6.3 |
| 2 | Fe xi | 182.17 | 6.2 | 18 | Fe xii | 196.64 | 6.3 | 34 | Fe xii | 312.25 | 6.3 |
| 3 | Fe x | 184.54 | 6.2 | 19 | Fe xiii | 200.02 | 6.3 | 35 | Mg viii | 313.74 | 6.0 |
| 4 | Ni xvi | 185.23 | 6.4 | 20 | Fe xiii | 202.04 | 6.3 | 36 | Si vii | 314.36 | 6.1 |
| 5 | Ca xiv | 186.61 | 6.6 | 21 | Fe xii | 203.83 | 6.3 | 37 | Mg viii | 315.02 | 6.0 |
| 6 | Fe xii | 186.89 | 6.3 | 22 | S xi | 205.82 | 6.3 | 38 | Si viii | 316.22 | 6.1 |
| 7 | Fe xii | 187.93 | 7.1 | 23 | Ni xvi | 288.17 | 6.4 | 39 | Mg viii | 317.03 | 6.0 |
| 8 | Fe x | 188.23 | 6.2 | 24 | Ni xvi | 291.98 | 6.8 | 40 | Fe xii | 318.13 | 6.3 |
| 9 | Fe x | 188.30 | 6.2 | 25 | Fe xii | 292.46 | 7.1 | 41 | Mg vii | 319.03 | 5.8 |
| 10 | Fe xiv | 192.03 | 7.2 | 26 | Si x | 291.76 | 6.2 | 42 | Si viii | 319.84 | 6.1 |
| 11 | Fe xii | 192.39 | 6.3 | 27 | Si x | 296.11 | 6.2 | 43 | Ni xviii | 320.57 | 6.8 |
| 12 | Fe xii | 192.83 | 6.2 | 28 | S xi | 291.54 | 6.3 | 44 | Fe xii | 320.81 | 6.3 |
| 13 | Ca xviii | 192.85 | 6.8 | 29 | Ca xviii | 302.19 | 7.0 | 45 | Fe xii | 323.41 | 6.3 |
| 14 | Fe xii | 193.51 | 6.3 | 30 | Fe xv | 302.33 | 6.3 | 46 | Fe xvii | 323.65 | 6.8 |
| 15 | Ca xviii | 193.87 | 6.6 | 31 | Fe xv | 304.89 | 6.3 | 47 | Fe xv | 327.03 | 6.3 |
| 16 | Fe xii | 195.12 | 6.3 | 32 | Fe xx | 309.29 | 7.0 | | | | |

In the U304 channel, the strongest lines are Fe xv 284.16 Å, blend S xii 288.42 Å+ Fe xiv 289.15 Å, Ni xviii 291.98 Å, selfblend Si ix 296.11+.21 Å, Ca xviii 302.19 Å, blend Fe xvii 304.82 Å+ Fe xv 304.89 Å, Mg viii 315.02 Å, and the blend Ni xviii 320.57 Å+ Fe xii 320.81 Å, which have intensities of order $2 \times 10^{-4}$ erg s cm$^{-2}$ and higher. The strongest line in the spectral region—the Si x/He ii blend with $\lambda \sim 304$ Å—was removed from the spectra before the analysis.

In the U304 channel, the strongest lines are Fe xv 284.16 Å, blend S xii 288.42 Å+ Fe xiv 289.15 Å, Ni xviii 291.98 Å, selfblend Si ix 296.11+.21 Å, Ca xviii 302.19 Å, blend Fe xvii 304.82 Å+ Fe xv 304.89 Å, Mg viii 315.02 Å, and the blend Ni xviii 320.57 Å+ Fe xii 320.81 Å, which have intensities of order $2 \times 10^{-4}$ erg s cm$^{-2}$ and higher. The strongest line in the spectral region—the Si x/He ii blend with $\lambda \sim 304$ Å—was removed from the spectra before the analysis.

Emission of hot spectral lines such as Fe xiv 192.03 Å ($T_m = 16$ MK), Ca xviii 192.85 Å ($T_m = 6.3$ MK), Fe xvii 292.46 Å ($T_m = 13$ MK), Ca xviii 302.19 Å ($T_m = 10$ MK), and Fe xx 309.29 Å ($T_m = 10$ MK) is produced only during flares. Spectral images of a flare in these lines are compact and usually not intermingled with other spectral lines (we
inspected a large number of the SPIRIT spectroheliograms). Thus, these lines can be used for detecting solar flare and they are ideal for high-temperature DEM and Doppler shift analysis.

The obtained spectra of all flares are similar, but still there are some differences. The absolute fluxes in separate spectral lines measured by SPIRIT in the M5.6 and X3.4 flares are similar and two-fold higher than those in the X1.3 and X17 flares. The decrease is in direct correlation with the decrease of the total flux in the EIT images. The decrease may be caused by variations in the solar irradiance—the M5.6 and X3.4 flares were registered at the end of 2001 (near the maximum of solar activity), the X1.3 flare was observed in 2004 July, and the X17 flare was observed in 2005 September (near the minimum of solar activity), as well as a degradation of EIT sensitivity (BenMoussa et al. 2013).

![V190 spectra of flares: X1.3 observed on 2004 July 16, M5.6 observed on 2001 September 16, and X17 observed on 2005 September 7. The black curve denotes the observational data, the blue vertical lines denote individual spectral lines from catalog, and the red curve denotes the fitted spectra (calculated using the instrumental FWHM).](image)

(A color version of this figure is available in the online journal.)

### 4.2. Plasma Diagnostics

The result of the DEM reconstruction is presented in Figure 9: red lines correspond to different runs (we used 100 runs), the black line is an average (median) DEM, and the green line denotes the initial DEM, obtained in the zero-level step.

The obtained DEMs have a similar shape—a local minimum at \( T \sim 0.6-0.8 \) MK (cold plasma), a local maximum at \( T \sim 2.5 \) MK (warm plasma), and a global maximum at \( T \sim 10 \) MK (hot plasma). The double-peaked shape may be associated with different structures: the warm plasma fill loops, which are adjacent to the flaring region (Schmelz et al. 2011), whereas the hot plasma is produced in the flaring region. The M5.6 and X1.3 flares have narrower hot-component peaks, which may be attributed to the earlier phases of the flare decays (\( \Delta T \sim 10 \) and 7 minutes after the flare maxima). The X3.4...
and X17 flares were registered on later phases (Δτ ~ 40 and ~3 hr after the flare maxima), therefore the hot plasma had time to warm up the surroundings. The warm-component peak in the latter two flares has a double-peaked shape with \( T_1 = 1.6 \text{ MK} \) (both flares) and \( T_2 = 4.0 \text{ MK} \) (the X17 flare). These double peaks in warm plasma may also be attributed to spatially separated structures.
The steep decrease in DEMs with $T > 10$ MK observed in the M5.6, X1.3, and X17 flares is determined by the intensities of hot lines, such as Ca xvii ($T_m = 6.3$ MK), Fe xxii ($T_m = 13$ MK), and Ca xviii ($T_m = 10$ MK), Fe xx ($T_m = 10$ MK). The primary contribution, however, is definitely due to the Fe xxiv 192.03 Å line, which has $T_m = 16$ MK. Since the V190 channel observations were unavailable for the X3.4 flare, it is possible that the DEM values with $T > 10$ MK are overestimated in the flare.

The confidence level of the DEMs is assessed by the relative spread of different DEM solutions and amounts to as much as a factor of two (each solution from the range equally well describes the observational data, so each solution from the range is equally possible).

The results of $n_e$ analysis are presented in Figure 10—the L-functions of the Fe xi, Fe xii, Fe xiii, and Fe xiv ions for the M5.6, X1.3, and X17 flares are given (black lines denote V190 and red lines denote U304). Whereas L-functions of a single ion should cross in a single point, one can see considerable disagreement in several cases. We note that during the iterative steps we tried to improve the agreement of the L-functions by varying the intensities of blended lines; however, better agreement was not reached.

The best consistency among different spectral lines is observed in the Fe xii ion. The $n_e$ values are $\sim 6.5 \times 10^9$ cm$^{-3}$ (the L-functions intersect in the range $\log n_e = [9.7–9.9]$) in all flares. The most reliable Fe xi lines—179.76, 180.41, 188.23 Å—also favor this value. The L-functions of the Fe xii ion show considerable discrepancy. The 200.02, 202.04, and the 203.83 Å lines cross at a density $n_e \sim 2 \times 10^9$ cm$^{-3}$, whereas the 196.54 and 200.02, and 203.80 + 0.83 Å lines cross at a density $\sim 6.5 \times 10^9$ cm$^{-3}$. The Fe xv lines show systematic discrepancies in all flares. We will discuss possible causes of the discrepancies in the next section. We used a value of $n_e = 6.5 \times 10^9$ cm$^{-3}$ for all the flares for DEM analysis and the calculation of synthetic spectra.

### Table 3

Intensities of Spectral Lines [$\times 10^{-5}$ erg s$^{-1}$ cm$^{-2}$] in the Solar Flares in the V190 Channel

| $N$ | Ion | $\lambda$ (Å) | $T_{\text{max}}$ (K) | Intensity | M5.6 | X1.3 | X3.4 | X17 |
|-----|-----|--------------|-----------------|---------|-----|-----|-----|-----|
| 1   | Ni xv | 176.10 | 6.4 | 3.6 ... | N/A ... |
| 2   | Ni xv | 176.74 | 6.4 | 6.8 6.4 | N/A ... |
| 3   | Ca xvii | 176.93 | 6.7 | 21.3 8.0 | N/A ... |
| 4   | Fe x | 177.24 | 6.1 | 44.6 38.8 | N/A 22.7 |
| 5   | Fe xi | 178.06 | 6.2 | 24.3 10.6 | N/A 4.9 |
| 6   | Ni xv | 179.27 | 6.4 | 11.8 8.9 | N/A 6.2 |
| 7   | Fe vi | 179.76 | 6.2 | 20.2 23.9 | N/A 16.2 |
| 8   | Fe vi | 180.41 | 6.2 | 180.89 | N/A 86.8 |
| 9   | Fe vi | 180.44 | 6.1 | 18.7 11.9 | N/A 13.4 |
| 10  | Fe xi | 180.60 | 6.2 | 12.9 6.9 | N/A 6.9 |
| 11  | Fe xi | 181.14 | 6.2 | 10.9 11.0 | N/A 6.7 |
| 12  | Fe xi | 182.17 | 6.2 | 35.7 25.6 | N/A 23.7 |
| 13  | Ca xiv | 183.46 | 6.6 | 6.1 6.6 | N/A 3.7 |
| 14  | Ni xiv | 183.97 | 6.3 | 33.6 ... | N/A ... |
| 15  | O vi | 184.12 | 6.3 | 3.4 13.4 | N/A 4.4 |
| 16  | Fe x | 184.54 | 6.1 | 35.6 30.7 | N/A 30.6 |
| 17  | Fe vi | 184.80 | 6.2 | 11.8 13.8 | N/A 8.0 |
| 18  | Fe vii | 185.21 | 5.8 | 11.5 11.3 | N/A 12.6 |
| 19  | Ni xvi | 185.23 | 6.4 | 21.3 12.0 | N/A 6.6 |
| 20  | Fe viii | 186.60 | 5.7 | 8.9 8.1 | N/A 14.0 |
| 21  | Ca xii | 186.61 | 6.6 | 13.0 11.5 | N/A 8.7 |
| 22  | Fe xii | 186.85 | 6.3 | 29.7 17.6 | N/A 15.0 |
| 23  | Fe xii | 186.89 | 6.3 | 62.5 36.6 | N/A 31.4 |
| 24  | Fe xii | 187.93 | 7.1 | 7.3 4.9 | N/A 5.6 |
| 25  | Ar xiv | 187.97 | 6.6 | ... ... | N/A 5.5 |
| 26  | Fe xii | 188.17 | 6.3 | 3.4 ... | N/A ... |
| 27  | Fe xii | 188.23 | 6.2 | 88.0 47.5 | N/A 43.8 |
| 28  | Fe xii | 188.30 | 6.2 | 33.0 17.8 | N/A 16.4 |
| 29  | Fe xii | 188.50 | 6.0 | 7.3 6.7 | N/A 10.0 |
| 30  | Si xii | 188.68 | 6.3 | 5.5 ... | N/A ... |
| 31  | Ar xi | 188.81 | 6.3 | 3.7 ... | N/A ... |
| 32  | Fe xii | 189.13 | 6.2 | 8.6 4.6 | N/A 4.6 |
| 33  | Fe xii | 189.72 | 6.2 | 6.6 3.6 | N/A 3.5 |
| 34  | Fe xii | 189.94 | 6.1 | 7.9 6.7 | N/A 5.0 |
| 35  | Fe xii | 190.04 | 6.2 | 14.4 9.2 | N/A 5.1 |
| 36  | Fe xii | 190.07 | 6.3 | 4.0 ... | N/A ... |
| 37  | Fe xii | 191.05 | 6.3 | 5.4 ... | N/A ... |
| 38  | Fe xii | 191.22 | 6.0 | ... ... | N/A 4.1 |

Notes. Columns correspond to different flares; rows denote different spectral lines. The flux in a particular spectral line corresponds to the whole flaring region (and does not contain a factor of 1/sr). Dots denote that the line was too weak in a particular spectrum and “N/A” in the X3.4 flare means that the flare was not observed in the V190 channel of SPIRIT.
4.3. Comparison of Observational and Theoretical Line Intensities

We compared the observational and theoretical intensities using different approaches: in the DEM reconstruction procedure, we used L-function plots and compared observational and synthetic spectra. All of these approaches, in essence, consist of comparing observational and theoretical line intensities, whereas each approach provides some additional convenience in terms of data analysis.

In the vast majority of spectral lines, the correspondence of observational and theoretical intensity is within factor of two. The most striking discrepancy $R$ (ratio of observational/theoretical intensity) is observed in the DEM reconstruction in the following lines: Fe xvii 323.41 Å (~8), Fe xiii 312.11 Å (~2.5), Fe xii 312.25 Å (~2.5), Fe xiii 202.04 Å (~2.0), Fe xvii 323.65 Å (~0.7), and Fe xv 327.05 Å (~0.7). There is also a systematic discrepancy in the relative intensities of the Mg viii and Si viii lines—whereas these spectral lines have a similar dependence on temperature and density, the ratio $R$ for Mg viii is constantly higher and for Si viii it is consistently smaller than 1. The L-function plots show discrepancies in the Fe xiii and Fe xv lines. Comparison of the observational and synthetic spectra reveal several discrepancies in other lines. The observed discrepancies are typical for the analyzed flares and we will discuss them all together.

The observed intensity of the Fe xii 323.41 Å line is approximately eight times higher than those predicted in the DEM reconstruction. The discrepancy cannot be attributed to problems with SPIRIT spectral sensitivity due to the good correspondence of other intense lines at nearby wavelengths. The Fe xii 323.41 Å line is blended with Fe xvii 323.65 Å; however, the spectral profile of the blend seems unlikely to be the blend of two close spectral lines. The incorrect identification of one of the lines seems quite reasonable.

The next two lines with ratio $R \sim 2.5$ are Fe xii 312.11 Å and Fe xii 312.25 Å. The lines fall within a wide blend, which encompasses Fe xii 311.55 Å, Ni xv 311.76 Å, Mg viii 311.76 Å,
Fe\textsc{xiii} and Fe\textsc{xii}, Co\textsc{xvii} 312.54 Å, Fe\textsc{xv} 312.56 Å, and Fe\textsc{xiii} 312.87 Å (the strongest lines according to the synthetic spectra). A detailed analysis of the blend deserves effort and attention; a quick look (using both spectra in Figure 8 and the L-functions in Figure 10) shows no simple solution for improving the $R$ ratios either via changes of $n_e$ or in changes of the relative intensity of lines involved in the blend. Fortunately, in the U304 channel there are a number of strong lines suitable for the reconstruction of intensities of the Mg\textsc{viii} and Fe\textsc{xiii} spectral lines.

The Fe\textsc{xiii} 202.04 Å line is among the most intense lines of the Fe\textsc{xiii} ion; however, its ratio $R$ is approximately 2 in all analyzed flares. The observed discrepancy cannot be attributed to problems with SPIRIT's spectral sensitivity, since the line falls between other strong Fe\textsc{xiii} lines—200.02 Å and a blend of 203.80 + 0.83 Å (the observed intensity of these lines is consistent with theory). During the DEM reconstruction, other Fe\textsc{xiii} lines (196.54, 200.02, 203.80, 312.11, 318.13, and 320.81 Å) were predicted with higher accuracy (usually better than 40%), eliminating possible issues with abundances or the temperature distribution of the emitting plasma. The L-functions of all Fe\textsc{xiii} lines are density sensitive (see Figure 10) and correction of the $n_e$ value seems reasonable and sufficient. However, the L-functions of the 200.02, 203.80, 203.83, and 320.81 Å lines have the same dependence on density and their absolute values are in a good agreement with theory. Density values $n_e$ obtained by the crossing of the 202.04 Å line and the 200.02 and 203.80 + 0.83 Å lines is systematically lower (~2 x $10^9$ cm$^{-3}$) than $n_e$ obtained with the Fe\textsc{xi} and Fe\textsc{xii} ions, which disfavors the 202.04 Å line. A good agreement with the density $n_e$, obtained with the Fe\textsc{xi} and Fe\textsc{xii} ions, was obtained by crossing the L-functions of the 196.54 Å line with the 200.02 and 203.80 + 0.83 Å lines. This result is in a slight contradiction with Brosius et al. (1998) and Shestov et al. (2009), who found good correspondence of the $n_e$ values measured by the Fe\textsc{xi} and Fe\textsc{xiii} lines. The L-functions of the 204.95, 312.11, and 321.40 Å lines show a similar behavior in density and in some flares their absolute values are in good agreement. However, the L-function of the 204.95 Å line does not produce a reasonable $n_e$ value (observed $n_e < 10^9$ cm$^{-3}$). In two flares (M5.6 and X1.3), the L-function of the 321.40 Å line has a common crossing with the 196.54, 200.02, and 203.80 + 0.83 Å lines, which confirms the correctness of the latter lines.

Figure 9. DEMs of the flares. Red lines denote DEMs obtained during different iterations of the GA procedure, black lines denote the average value, and the green line denotes the DEM, obtained during zero-level analysis.

(A color version of this figure is available in the online journal.)
Figure 10. L-function of Fe XI–Fe XIII and Fe XV ions for the M5.6, X1.3, and X17 flares. Black lines denote spectral lines from the V190 channel; red lines denote the U304 channel.

(A color version of this figure is available in the online journal.)
Given the above information, the overall agreement of the Fe xiii L-functions may be improved by decreasing the L-functions of the 202.04 Å line by a factor of ~2 and the 204.95 Å line by a factor of ~4. The observed excess in the L-functions may be caused by unaccounted blending of spectral lines or inappropriate atomic data. The atomic structure of the Fe xiii ion has recently been extensively studied by Del Zanna (2011) and the author did not find any problems with the ion. However, the observed discrepancies in the SPIRIT data (with no strong blend candidates provided by CHIANTI) and the inconsistency in the $n_e$ values obtained with Fe xiii indicate that some questions still remain.

The Fe xvii 323.65 Å line has a ratio $R \sim 0.7$ and is blended with Fe xii 323.41 Å. The latter line shows a striking discrepancy with $R \sim 8$. The spectral profile of the blend seems unlikely to be the blend of two close spectral lines and an incorrect identification of one of the lines seems quite reasonable.

The Fe xv 327.03 Å line has a ratio $R \sim 0.7$. It is blended with K xvii 326.78 Å the ratio $R$ is not too bad taking into account that both lines are on the edge of the SPIRIT spectral range.

The other issue is the systematic discrepancy of intensities between the Mg vii and Si viii spectral lines. The two ions have similar atomic structure ($2s^22p^2 \rightarrow 2s^22p^3$ transitions in Mg vii and $2s2p^n \rightarrow 2s2p^{n+1}$ transitions in Si viii), similar abundances (both in coronal and photospheric models), close wavelengths—311.77, 313.74, 315.02, and 317.03 Å (Mg vii) and 314.36, 316.22, and 319.84 Å (Si viii)—and the contribution functions of the lines have a similar dependence on temperature and density. Nevertheless, the ratios $R$ for the Mg vii lines are 1.6, 1.4, and 1.03 for the 313.74, 315.02, and 317.03 Å lines (averaged by flares), whereas the ratios $R$ for Si viii approach 0.7, 0.6, and 0.8 for the 314.36, 316.22, and 319.84 Å lines, respectively. Inadequate abundances are the most likely cause of the discrepancy. A similar possibility was pointed out by Schmelz et al. (2012) in their analysis of SERTS data.

The L-functions of the Fe xv ion can be separated into two groups: those that decrease with density (the 304.89, 307.75, and 321.77 Å lines) and those that do not change with density (the 284.16, 312.56, 324.98, and 327.03 Å lines). The L-functions inside each group should agree. According to the observational data, the L-function of the 284.16 Å line is usually four times lower than the others. The discrepancy can be caused by different factors: the SPIRIT detector saturation, problems with SPIRIT instrument calibration, the uncalibrated SPIRIT data. Using spectroscopic instrumentation with relatively high spectral resolution could calibrate the uncalibrated SPIRIT data. The procedure does not take into account unknown lines and this is the main disadvantage of the proposed method of the identification. However, a comparison of the observed and calculated intensities provides a lot of information about the reliability (qualitative—correct or incorrect) and accuracy (quantitative—say 10% or 50%) of the identification. The accuracy depends on two factors—(1) the quality of the observational data (the accuracy of the spectral calibration, low noise, absence of scattered light in the instrument, the compactness of the emitting plasma, etc.) and (2) the accuracy of the atomic data used.

We analyzed how the identification of the obtained spectra depended on the relative calibration of SPIRIT and concluded that the relative calibration is better than a factor of two (there are a number of lines whose intensity ratio does not depend on density and that agree with theory). The absolute calibration was obtained using simultaneous EIT images—195 Å for the V190 channel and 304 Å for the U304 channel. After the absolute calibration, the average line intensities in the both channels (which were calibrated independently) agree with each other. However, the resulting absolute calibration of SPIRIT is as good as the EIT calibration. Any errors in EIT calibration—for example, due to the decay of the EIT sensitivity (BenMoussa et al. 2013)—will affect the absolute calibration of the presented spectra.

We analyzed how the results of the DEM reconstruction depend on the errors (up to 30%), artificially introduced into the observational data. The obtained discrepancy turned out to be within the DEM confidence level, obtained in the GA. Other factors (beside adequate identification) do not play such an important role in the final accuracy. We assess the final accuracy of the observational data (including the absolute calibration) to be a factor of four.

The analysis of simultaneous EIT and SPIRIT data proved that telescopic and spectroscopic observations significantly enhance each other. SPIRIT gives direct information about the relative flux in each spectral line contributing to the EIT image and the relative flux measured by EIT (in dn units) allows us to calibrate the uncalibrated SPIRIT data. Using spectroscopic instrumentation with relatively high spectral resolution could enhance the informational content of other instruments, like the Atmospheric Imaging Assembly (AIA) or EVE.

The other important aspect for spectroscopic analysis is spectral resolution. The instrumental resolution of SPIRIT is $\sigma \sim 0.1$ Å (for comparison, the spectral resolution of EIS is $\sigma \sim 0.020$ Å). Many strong and important lines are not resolved by SPIRIT, such as the Fe xiii 203.80 and 203.83 Å lines, or the Ni xvi 291.98 Å, Fe xv 292.26 Å, or Fe xii 292.45 Å lines. Nevertheless, the identification procedure used allowed us to deconvolve the blends and calculate the intensities. In some cases, we obtained a good correspondence between the observational and theoretical intensities: an example is the blend of the Fe xiii 203.80 and 203.83 Å lines, which are in a good agreement with the Fe xiii 200.02 Å line; the blend of the Fe xiii 196.54 Å and Fe xii 196.64 Å lines also agrees with the rest of the Fe xii and Fe xiii lines.

5. DISCUSSION

During the identification, we fit the observational data with synthetic spectra calculated using CHIANTI. The procedure does not take into account unknown lines and this is the main disadvantage of the proposed method of the identification. However, a comparison of the observed and calculated intensities provides a lot of information about the reliability (qualitative—correct or incorrect) and accuracy (quantitative—say 10% or 50%) of the identification. The accuracy depends on two factors—(1) the quality of the observational data (the accuracy of the spectral calibration, low noise, absence of scattered light in the instrument, the compactness of the emitting plasma, etc.) and (2) the accuracy of the atomic data used.

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Figure 11. Comparison of observational and synthetic (calculated using DEM and \( n_e \) with CHIANTI) spectra for the M5.6 flare. Top panel: V190 channel. Bottom panel: U304 channel. The black curve denotes the observational data, the blue vertical lines denote individual spectral lines, and the red curve denotes synthetic spectra (calculated using the instrumental FWHM). The arrows show major the discrepancies between the observational and synthetic spectra.

(A color version of this figure is available in the online journal.)

was not achieved: the multiple blend with \( \lambda \sim 312 \, \text{Å} \) is an example, where the lines contributing to the blend show a poor correspondence with the theory. Nevertheless, the method of deconvolving blends using a synthetic spectrum is a powerful tool for spectroscopic analysis.

Let us now compare the calculated DEMs with those obtained from other instruments. The DEM that is widely used for modeling EUV spectra is that presented by Dere & Cook (1979); this DEM is actually provided by CHIANTI as a flare.dem. The authors analyzed the decay phase of an M2 flare using observations from the S082A EUV spectroheliograph and the S082B UV spectrograph onboard Skylab. During the DEM calculation, the authors used the quantity “total line power radiated by the plasma” (which actually agrees with our approach). However, the DEM values provided by CHIANTI are expressed in units of \( \text{cm}^{-5} \, \text{K}^{-1} \). The DEMs, calculated in our analysis correspond to the whole flaring region and we need to assess the area associated with the flare. In order to assess the area, we inspected monochromatic images of the flares and we need to assess the area associated with the flare. In order to assess the area, we inspected monochromatic images of the flares and concluded that the images have a symmetrical Gaussian shape with a typical FWHM \( \sim 8 \) pixels (good examples are the bright lines of the Ni \( \text{xviii} \), S \( \text{xi} \), and Ca \( \text{xviii} \) ions in Figure 3). This is the minimum size of the structure observed on the spectroheliograms and the size is determined by the point spread function of the instrument (primarily, its focusing mirror). Thus, the spatial size of a flare should not exceed \( \sim 3 \times 3 \) pixels so as not to increase the flare images. Taking into account the angular size of a pixel of SPIRIT—6.7 arcseconds—we obtain that \( 3 \times 3 \) pixels corresponds to the area \( 2 \times 10^{18} \, \text{cm}^2 \).

We multiply flare.dem from CHIANTI by this factor and compare it with the current DEMs. The correspondence is good enough: beside a similar shape (the local minima and maxima agree in the two data sets), we also find an order-of-magnitude agreement—\( \sim 8 \times 10^{40} \, \text{cm}^{-3} \, \text{K}^{-1} \) (cold component) and \( \sim 6 \times 10^{42} \, \text{cm}^{-3} \, \text{K}^{-1} \) (hot component).

We performed another verification of the calculated DEMs: we simulated GOES X-ray fluxes using the calculated DEMs and GOES response functions (goes_res2.dat from Solar Software). The calculated fluxes agreed within an order of magnitude or better with those fluxes actually measured by GOES.

To investigate heating dynamics in observed flares—impulsive or continuous—we compared the hot component lifetime (\( \tau_{\text{life}} \), a few hours) with its conductive cooling time (\( \tau_{\text{cond}} \)): if \( \tau_{\text{cond}} \ll \tau_{\text{life}} \), then heating is continuous; if \( \tau_{\text{cond}} \sim \tau_{\text{life}} \), then heating is impulsive. We estimated \( \tau_{\text{cond}} \) with the formula (Culhane et al. 1994):

\[
\tau_{\text{cond}} = \frac{21 n_e k_b h^2}{5 \kappa T^{5/2}},
\]
where \( n_e \) is the plasma electron density, \( \kappa \) is Boltzmann’s constant, \( h = 1.5 \times 10^7 \text{ m} \) is the characteristic size, \( \kappa = 9.2 \times 10^{-7} \text{ erg s}^{-1} \text{ cm}^{-1} \text{ K}^{-7/2} \) is the Spitzer conductivity, and \( T = 10 \text{ MK} \) is the plasma temperature. The electron density in flares ranges from \( 10^{10} \text{ cm}^{-3} \) to \( 10^{12} \text{ cm}^{-3} \) (Milligan et al. 2012). For \( n_e = 10^{10} \text{ cm}^{-3} \), \( \tau_{\text{cond}} \sim 30 \text{ s} \) (\( \tau_{\text{cond}} \ll \tau_{\text{life}} \)), which requires continuous heating; for \( n_e = 10^{12} \text{ cm}^{-3} \), \( \tau_{\text{cond}} \sim 1 \text{ hr} \) (\( \tau_{\text{cond}} \sim \tau_{\text{life}} \)), which favors impulsive heating. SPIRIT spectra do not have high-temperature spectral lines suitable for the plasma DEM and \( n_e \), on calculation of synthetic spectra and the measurement of intensity for as many as 70 spectral lines and were registered on different stages of their decay, the registered spectra, as well as the proposed identification of abundances or atomic rates, etc. which may involve, along with a refinement of the spectral calibration and line identification, so we estimate \( n_e \) using the obtained DEMs:

\[
EM_{\text{hot}} \sim n_e^2 h^3 \Rightarrow n_e \sim \sqrt{EM_{\text{hot}}/h^3} \sim 10^{11} \text{cm}^{-3}. \tag{5}
\]

This is a rough estimate: the accuracy of the DEMs is a factor of four and the real volume of the hot component is probably less than \( h^3 \) (most likely it is not spherical, but has a loop-like geometry). So, \( n_e \) is probably closer to \( 10^{12} \text{ cm}^{-3} \) than \( 10^{10} \text{ cm}^{-3} \) and heating in the observed flares is most likely impulsive.

6. CONCLUSIONS

Initially, the main goal of the work was to present unique observational data—EUV spectra of large solar flares, observed by the SPIRIT spectroheliograph. Due to the relatively low spectral resolution of SPIRIT, many lines are blended, which prevents a straightforward method for line identification and measurement. The original procedure for spectra analysis, based on calculation of synthetic spectra and the measurement of plasma DEM and \( n_e \), not only allowed for the identification and measurement of intensity for as many as 70 spectral lines in each spectral band in each flare, but also provided a lot of other important information. The spectroscopic analysis demonstrated the accuracy of the adopted spectral calibration of the SPIRIT spectroheliograph. Simultaneous observations of the EIT telescope and the SPIRIT spectroheliograph allowed the calculation of absolute fluxes in each spectral line.

Whereas the analyzed flares belong to different X-ray classes and were registered on different stages of their decay, the registered spectra and the calculated DEMs have many properties in common. All DEMs have a similar shape with global maxima in each spectral band in each flare, but also provided a lot of other important information. The spectroscopic analysis demonstrated the accuracy of the adopted spectral calibration of the SPIRIT spectroheliograph. Simultaneous observations of the EIT telescope and the SPIRIT spectroheliograph allowed the calculation of absolute fluxes in each spectral line.

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