Data Article

Data on single pulse fs laser induced submicron bubbles in the subsurface region of soda-lime glass

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ABSTRACT

Submicron bubble formation in the subsurface range of soda-lime glass is investigated. The bubbles are induced by single femtosecond laser pulse irradiation with the wavelength of \( \lambda = 775 \text{ nm} \), the pulse duration of \( t_p = 150 \text{ fs} \) and the laser beam diameter of \( \sim 12 \mu\text{m} \). The data shows the changes of the morphologies of the soda-lime glass after laser irradiation with different pulse energy. Moreover, the data shows the detail of the cross-section view of the bubble during the Focused ion beam (FIB) cutting. It is found that the bubbles can be formed in a rather narrow pulse energy range with the bubbles in the size of \( 300 \text{ nm} - 3 \mu\text{m} \) which is much smaller than the laser beam diameter. Data presented in this article are related to the research article “Submicron bubbles/voids formation in the subsurface region of soda-lime glass by single pulse fs laser-induced spallation” [1].

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1. Data description

The morphologies of the soda-lime glass after laser irradiation with different laser pulse energy are shown in Fig. 1 with optical microscopy (OM) measurements. The bubbles are formed with a moving laser source while the scanning speed is 4 mm/s and the repetition rate of the laser is 1 kHz. The three kinds of morphologies characterize the three different surface conditions induced by laser irradiation: a) surface modification; b) bubble formation and d) laser ablation crater.

Table 1 shows the height or the depth of the structures, the morphologies and the profile of the measured bubbles/crater in dependence on the applied pulse energy. The sizes of the bubbles are measured by the atomic force microscopy (AFM) as well as the white light interference microscopy (WLIM). The morphologies are measured by the scanning electron microscopy (SEM). The bubbles can be obtained with the maximum height of ~500 nm in the range of the surface modification and ablation. The pulse energy that below 5.32 μJ (F = 4.01 J/cm²) can only induced the surface modification on the soda-lime glass or some small bubbles with the height below 60 nm. The glass features small bulging to the direction of the incident laser beam but the height is only several nm. When the pulse energy is higher than 5.5 μJ (F = 4.14 J/cm²), most of the laser irradiated spots show explored bubbles or ablation craters. The size of the bubbles or the ablation craters are much smaller than the laser beam diameter. From the table it can be found that with the increasing of the incident pulse energy, the bubbles’ height increases too as well as the bubble diameter. But there are some fluctuations of the bubble size due to the surface condition of the soda-lime glass.

Fig. 2 shows the FIB cutting process of the bubble and the cross-section view of the bubble. A thin Pt film ~1.5 μm is deposited on the glass surface in order to protect the surface structures as well as the
bubble. The bubble is fabricated with 5.32 \( \mu \)J (\( F = 4.01 \) J/cm\(^2\)). From the cross-section view of the bubble it can be clearly seen that the bubble inside is a hollow ellipsoid with an upper shell of about 100 nm in thick. This thickness shows constant in relation to the other bubble shell with damaged or flipped bubbles presented in Ref. [1].

Fig. 3 sketches the laser irradiation setup as well as the measurement of the laser signal by the oscilloscope. The oscilloscope is connected to the photodiode which is fixed near the beam path to detect the scatter light. Fig. 4 shows the probability of the bubble formation at a certain pulse energy range from 5.5 \( \mu \)J to 5.6 \( \mu \)J.

## 2. Experimental design, materials, and methods

### 2.1. Experimental design

A linear polarized femtosecond laser with a wavelength of 775 nm, a pulse duration of 150 fs, and a pulse repetition of 1 kHz is used in this study. The laser beam with a Gaussian profile is focused on the surface of the glass with a single lens having a focal length of 60 mm. The laser spot radius was determined to be \( r_g = 6.5 \pm 0.5 \) \( \mu \)m. The samples are irradiated in a pulse energy range of \( E = 3.00 \) (\( F = 2.26 \) J/cm\(^2\)) \( \sim \) 20 \( \mu \)J (\( F = 15.1 \) J/cm\(^2\)) with a scanning speed of 4 mm/s and then irradiated with single laser pulses arranged in a line as shown in Fig. 1. After that, a spacing of more than 20 \( \mu \)m is used for single pulse irradiation to avoid the influence of the overlap by multi-pulse irradiation with the laser pulse energy varied from 4.6 (\( F = 3.47 \) J/cm\(^2\)) to 6.6 \( \mu \)J (\( F = 4.97 \) J/cm\(^2\)).
Table 1
The formed three kinds of surface morphologies including modification, bubbles formation and ablation independence on the incident pulse energy and their relative measurements including pulse signal measured by the oscilloscope (OSC), the profile and 3D topography measured by AFM, WLIM, SEM, OM and the related parameters.

| Number | Surface modification | Small bubble | Increasing bubble size with increasing pulse energy |
|--------|----------------------|--------------|---------------------------------------------------|
| OSC    | ![Image](image1.png)  | ![Image](image2.png) | ![Image](image3.png) |
| AFM    | ![Image](image4.png)  | ![Image](image5.png) | ![Image](image6.png) |
| Profile| ![Image](image7.png)  | ![Image](image8.png) | ![Image](image9.png) |
| WLIM   | ![Image](image10.png) | ![Image](image11.png) | ![Image](image12.png) |
| Profile | WL | IM | SE | OM |
|---------|----|----|----|----|
|         | ![Image](image1) | ![Image](image2) | ![Image](image3) | ![Image](image4) |
|         | ![Image](image5) | ![Image](image6) | ![Image](image7) | ![Image](image8) |
|         | ![Image](image9) | ![Image](image10) | ![Image](image11) | ![Image](image12) |
|         | ![Image](image13) | ![Image](image14) | ![Image](image15) | ![Image](image16) |
|         | ![Image](image17) | ![Image](image18) | ![Image](image19) | ![Image](image20) |
| E       | 5.44 μJ | 4.10 J/cm² | -180 nm |
| F       | 5.388 μJ | 4.06 J/cm² | 385 nm |
| H       | 5.369 μJ | 4.04 J/cm² | 308 nm |
|         | 5.334 μJ | 4.02 J/cm² | 420 nm |
| Parame  |     |     |     |     |
| ter     |     |     |     |     |
2.2. Materials

The material used in this experiment is microscope slides (extra white soda-lime glass, Menzel Thermo Scientific) with a thickness of 1 ± 0.1 mm. The composition of the glass is 72% SiO₂, 14% Na₂O, 6% CaO, 4% MgO, 1% Al₂O₃, 1% K₂O and further additives in small amounts. The samples are cleaned in an ultrasonic bath with isopropanol and dried by nitrogen gas flow. However, each piece of the sample will have some difference on the composition based on the fabrication process which may result in the bubble size deviation in Table 1.

Fig. 2. The FIB cutting process of the bubble, the arrows in the image shows the border of the bubble from the upper shell of the bubble to the platinum film and the border of the glass substrate to the platinum film.

Fig. 3. Sketch of experimental setup for single pulse energy measurement and the equivalent circuit for detecting the laser signal and the measured voltage signal in correlation of the laser pulse energy.
2.3. Methods

In the experimental design it has been mentioned that arrays of laser spots have been built with single pulse laser irradiation ordered in rows of nominal equal irradiations and varying the fluence from row to row. At this condition, a photodiode (THORLABS, DET 10 A/M, 200–1100 nm) is fixed near the beam path to measure the scatter light signal during the laser irradiation for each laser pulse. An oscilloscope is used to record the voltage from the photodiode so that the single pulse laser energy can be recorded. Due the extremely short pulse duration (150 fs) and the inductive and capacitive properties of the measuring circuit, the electrically recorded pulse is not proportional to the temporal pulse power but proportional to the pulse energy. As the impedance of the devices can be considered as constant, the integral value of the voltage measured by the oscilloscope provides the laser pulse energy. The equivalent circuit is shown in Fig. 3.

Before the laser irradiation process, an energy meter (Coherent) is placed in the laser beam path between the mirror and the focus lens. The pulse energy is measured by the energy meter while the photodiode connected to the oscilloscope detects the scatter light at the same time. After that we calculate the average pulse energy of the scores of the pulse energy $E_a$ as well as the relative integral value $I_a$ detected from the oscilloscope. We define $k = E_a/I_a$ for normalization so that we can get the scatter light signal and its integral value $I_i$. Therefore, $E_i = k \cdot I_i$. Therefore, the relationship between the pulse energy and the scatter light detected by the photodiode is established.

From Fig. 1 in row B it can be found that at the bubble formation range, there still have some fluctuations of the pulse energy which result in the morphology fluctuation and make the bubbles’ formation occur in a random probability. From the hundreds spots the probability of the formation of the different surface morphologies (checked by SEM) in dependence on the laser pulse energy has been summarized in Fig. 4. The bubbles appear in the energy range of 5.20 ($F = 3.99$ J/cm$^2$) – 5.30 μJ ($F = 3.99$ J/cm$^2$) but only with a rather low probability less than 10% but almost all the surface conditions feature the modification. However, when the pulse energy increases to 5.3 ($F = 3.99$ J/cm$^2$) – 5.4 μJ ($F = 4.07$ J/cm$^2$), there is an almost 50% probability with the bubble formation. Moreover, in the range of 5.35 ($F = 4.03$ J/cm$^2$) – 5.40 μJ ($F = 4.07$ J/cm$^2$), more than 60% probability can the laser induced bubbles. Therefore, we take the nearly energy range that can generate the bubbles with 50% probability.
to be the threshold of the bubbles’ formation, and in the range of 5.35 \( (F = 4.03 \text{ J/cm}^2) \) ~5.40 \( \mu \text{J} \) \( (F = 4.07 \text{ J/cm}^2) \), the bubbles can be generated stably. When the energy increases, the probability of the bubble formation decreases while the ablation probability increases.

The bubbles are induced by single fs laser pulse, so that the incubation effect can be neglect in this condition while the multiphoton absorption process must be considered during the laser irradiation. From the experimental results it can be found that the bubbles are all smaller than the Gaussian beam radius. The multiphoton absorption induced bubbles’ size \( r_{MP} \) can be estimated theoretically according to function (1) based on the Gaussian distribution when the incident energy threshold is higher than the formation threshold \([2]\). Here, \( r_R \) is the actual Gaussian beam diameter and \( n \) is the photon number for the multiphoton absorption. Therefore, for different number of photon absorption such as 2, 3 and 4, the effective interaction diameter can be estimated for 2, 1.6 and 1.4 \( \mu \text{m} \) which is in correlation with the bubble size shown in Table 1.

\[
r_{MP} = \frac{r_R}{\sqrt{n}}
\]  

(1)

The morphological measurements are performed using scanning electron microscope (SEM) and atomic force microscope (AFM). The composition in the near surface of the bubble is measured by SIMS and EDX.

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**Conflicts of Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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