The influence of CeF₃ on radiation hardness and luminescence properties of Gd₂O₃–B₂O₃ glass scintillator

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The effect of CeF₃ concentration and γ-irradiation on the physical, optical and luminescence properties of Gd₂O₃–B₂O₃–CeF₃ glasses were studied in this work. Before irradiation, the addition of CeF₃ in glass degraded the network connectivity observed from FTIR and possibly created the non-bridging oxygen (NBO) in glass structure. This NBO caused the reduction of Ce³⁺/Ce⁴⁺ ratio in XANES, the red-shift in transmission spectra and the raise of refractive index with addition of CeF₃ content. Such red-shift also was influenced by 4f–5d transition of Ce³⁺ dopant. This ion generated the strongest photoluminescence (PL) and radioluminescence (RL) in 0.3 mol% CeF₃-doped glass with nanoseconds decay time. The irradiation with γ-rays damaged the glass structure, broke the chemical bonds, and created color center in the glass network. The non-bridging oxygen hole center (NBOHC), that absorbed photons in the visible light region, caused the darkening, color change and increment of refractive index. These irradiation effects on glass were mitigated by the addition of CeF₃ that the electron donation of Ce³⁺ decreased the number of NBOHC. The Ce³⁺/Ce⁴⁺ ratio in most glasses after irradiation then reduced compared to them before irradiation, resulting to the decrease in PL and RL intensity. Our results confirm that CeF₃ can enhance the radiation hardness of glass and the 0.3 mol% CeF₃-doped glass is a promising glass scintillator.

Single crystal scintillators are used in various applications such as medical imaging, non-destructive inspection, nuclear or high energy physics, environmental monitoring and geological exploration. In radiation detectors, single crystals offer the advantage of having high light yields and fast response times. However, single crystal growth is an expensive and slow process; and single crystals can only be produced with limited shapes and sizes. On the other hand, glasses with various shapes and sizes are cheaper and faster to fabricate. Recently, the glass scintillators have been developed and several works have shown sufficient high light yields and fast decay times for practical applications, including the interaction of radiation with glass and their shielding properties. Investigation of novel glasses for radiation detection is therefore emerging, with particular focus on understanding the irradiation effects and improving the radiation hardness. The radiation hardness is the resistant of material that its properties was not changed or distorted by irradiation.

Gd₂O₃–B₂O₃-based glasses are suitable scintillators owing to their radiation interaction. The ¹⁰B boron isotope possesses a high capture cross-section for thermal neutrons, making it a suitable neutron detector. Additionally, B₂O₃ host glass is highly transparent, with good physical and chemical properties that meet the requirements for a scintillator. The high phonon energy of borate glass decreases its luminescence efficiency, but this can be mitigated by adding a heavy metal oxide, such as Gd₂O₃, into the glass. For γ-rays and X-rays detection, the addition of Gd increases the glass density and effective atomic number which improves the interaction between glass and such incoming radiation and the Gd³⁺ ion can efficiently transfer the energy to luminescence.

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centers such as lanthanide ions (Ln³⁺)²⁰,³⁰,³¹. In case of neutron detection, the ¹⁵⁵Gd and ¹⁵⁷Gd isotopes own a high capture cross-section for thermal neutrons²⁰,²⁷,²⁸. However, there is low energy γ-rays emitted from Gd under neutron irradiation, and distinguish this γ-rays from background γ-rays is technically impossible in the pulse height or spectroscopy-based techniques. Therefore, there is no one uses Gd-based scintillators for actual applications except for some special applications. If consider in the common pulse shape discrimination it may be possible to use³³. Therefore, the Gd₂O₃–B₂O₃-based glasses is very attractive for γ-rays and X-rays scintillator, but it has a difficulty for using in neutron detection.

Among the Ln³⁺ ions, trivalent cerium (Ce³⁺) is the most favored luminescence center for scintillator applications because the 5d–4f dipole allowed transition in Ce³⁺ results to a bright luminescence emission with nanoseconds decay time²⁴–²⁶. Previous works have investigated the Ce³⁺-doped Gd₂O₃–B₂O₃-based glasses such as Ce³⁺: Li₂O–B₂O₃–Gd₂O₃–B₂O₃, Ce³⁺: Gd₂O₃–CaO–SiO₂–B₂O₃, Ce³⁺: Li₂O–Gd₂O₃–BaO–B₂O₃, Ce³⁺–Dy³⁺: CaCO₃–ZnO–Gd₂O₃–B₂O₃, and Ce³⁺: Gd₂O₃–B₂O₃. In particular, our work on the xCeF₃-doped pure binary 27.5Gd₂O₃–(72.5–x)B₂O₃ (Ce:GB)³⁹ demonstrated the significant progress in binary glass preparation as the glass sample was successfully synthesized without adding any glass modifier compound to help in glass melting process. Our technique has the advantage of excluding unnecessary oxide components which possibly degrade the color, optical and luminescence properties of the glass. Consequently, the Ce:GB glass exhibited characteristics that make it a promising glass scintillator.

In order to fully capitalize on the potential of Ce:GB and other glasses as scintillators, in-depth investigation about the effects of irradiation on the glass’ properties, especially on its luminescence properties and radiation-hardness, are necessary for further study. Consequently, the Ce:GB glasses were irradiated by gamma-rays (γ-rays) and the effect of this irradiation on glass properties were invein binary glass preparation as the glassstigated in this present work.

Methods

Ce:GB glasses with 27.5Gd₂O₃–(72.5–x)B₂O₃–xCeF₃ composition were synthesized by a melt-quenching technique. The glass samples, 0CeGB, 0.05CeGB, 0.1CeGB, 0.3CeGB, 0.5CeGB, 1.0CeGB and 1.5CeGB contain different CeF₃ concentrations, with x being 0.00, 0.05, 0.10, 0.30, 0.50, 1.00 and 1.50 mol%, respectively. Details of the raw chemicals and the glass preparation procedure are stated in our previous work³¹. The Ce:GB glasses were irradiated with γ-rays carrying 1.17 and 1.33 MeV energies from a cobalt-60 (⁶⁰Co) source. The ⁶⁰Co source was calibrated with water and had a dose rate of 36.82 Gy/h at a distance of 1.0 m. The samples were placed at approximately 10 cm away from the radiation source for 6 h. The irradiation was performed at room temperature and in ambient atmosphere. The estimated irradiation dose rate and total dose on samples are 0.57 kGy/hour and 3.44 kGy, respectively. The oxidation state of Ce ion dopant in glasses were monitored by X-ray absorption near edge structure (XANES) spectroscopy at the Synchrotron Light Research Institute (SLRI), Thailand. The glasses densities (ρ) were determined using a 4-digit microbalance (AND, HR-200) and the Archimedes’ method with deionized water as an immersion liquid. The molar volumes of the glasses (Vₘ) were calculated using the relation: Vₘ = Mₑ/ρ. The refractive indices (n) of the glasses were measured by Abbe refractometer (Atago, DR-M2/M4) using the D-line (589 nm) source and 1-bromonapthalene as the contact liquid. Fourier-transform infrared (FTIR) spectra were recorded using an FTIR spectrometer (Agiilent, Cary 630). An Ultraviolet–Visible–near infrared (UV–VIS–NIR) spectrophotometer (Shimadzu, UV-3600) was used to measure the transmittance spectra. The phospholuminescence (PL) spectra of glasses were monitored by a spectrophotometer (Agiilent, Cary Eclipse) with xenon lamp as a light source. The PL decay profiles were obtained using the third harmonics (3ω, 290 nm) of a Ti:sapphire laser. The decay times were measured using a 25 cm focal length spectrograph which was fitted with a 600 grooves mm⁻¹ grating that was coupled to a Hamamatsu C1587 streak camera unit and a charge-coupled device (CCD) camera. For the X-ray induced optical luminescence or radioluminescence (RL) spectra, the glass samples were excited by X-rays from a Cu target generator (Inel, XRGAD-E) with 50 kV and 30 mA power. The RL emission signal was detected by an optical fiber and a spectrometer (Ocean Optics, QE65 Pro).

Results

The glass appearance, density and molar volume. Photographs showing the physical appearance of Ce:GB glasses before and after γ-ray irradiation are represented in Fig. 1. Before irradiation, the CeF₃-free glass (0CeGB) was highly transparent and colorless; while the color of the CeF₃-doped glasses (0.05CeGB, 0.1CeGB, 0.3CeGB, 0.5CeGB, 1.0CeGB and 1.5CeGB) became more greenish yellow as the amount of CeF₃ increased. After irradiation, the 0CeGB glass was dramatically darkened and least transparent, indicating that there was significant damage from the γ-rays. On the other hand, the irradiated CeF₃-doped glasses was less darkened and hence more transparent than the 0CeGB glass. A greenish yellowing in the glasses can be observed which the 1.5CeGB glass visually exhibited a similar level of transparency and color tone before and after irradiation.

Table 1 shows the density (ρ) and molar volume (Vₘ) of all Ce:GB glasses before and after γ-irradiation. The density of Ce:GB glasses were quite high in a range of 4.09–4.16 g/cm³, which are suitable for radiation detection. The CeF₃ concentration and irradiation did not seem to significantly affect the density and molar volume of glasses. Generally, the density of glass decreases if glass is irradiated by huge γ-rays that ejects the anions in structure. However, that density change is very small which requires high accurate measuring system. In order to observe and the irradiation dose using in this work was not high. The increment of density after irradiation in 0.05CeGB, 0.1CeGB and 0.3CeGB glasses were lower than 1%, so they could be the typical errors from measurement.

The oxidation state of cerium ion in glass. The typical XANES spectra represent the Ce L₂₃ edge of 0.05CeGB, 0.3CeGB and 1.5CeGB glasses before (Fig. 2a) and after irradiation (Fig. 2b), compared to the unir-
radiated standard compounds, CeF₃ and CeO₂. The XANES spectra show that the +3-oxidation state of Ce³⁺ in CeF₃ compound has a prominent absorption peak at 5727 eV, while the +4-oxidation state of Ce⁴⁺ in CeO₂ powder has obvious double peaks at 5731 eV and 5738 eV. By comparing both standard compounds, there was also a weak absorption peak of Ce³⁺ in the CeO₂ powder. Likewise, there were weak peaks of Ce⁴⁺ in the CeF₃ compound. These indicate that the cerium ions in CeF₃ and CeO₂ coexisted in both Ce³⁺ and Ce⁴⁺ states. The Ce:GB glasses in this work were doped with CeF₃. Therefore, their XANES spectra mimicked the spectrum of CeF₃ standard where the Ce³⁺ ion is dominant. The XANES data were evaluated using the Athena software to ascertain the quantity percentage of Ce³⁺ and Ce⁴⁺ ions in Ce:GB glasses. Before irradiation, the ratio of Ce³⁺/Ce⁴⁺ in the glasses decreased as the CeF₃ concentration increased. The same trend was also observed in CaO–SiO₂–B₂O₃–CeF₃ and SiO₂–Al₂O₃–Li₂O–Na₂O–K₂O–BaO–SrO–Tb₂O₃–Gd₂O₃–CeO₂ glass fabricated in air atmosphere by Rajaramakrishna et al. and Zu et al., respectively. On the other hand, the Ce³⁺/Ce⁴⁺ ratio in the glasses after irradiation increased with the increase in CeF₃ concentration. Considering the effect of γ-irradiation, it decreased the Ce³⁺/Ce⁴⁺ ratio of 0.05CeGB and 0.3CeGB glasses while it slightly increased this ratio in 1.5CeGB glass.

The glass network. Results of the FTIR measurements for the 0CeGB and 1.5CeGB glasses before and after γ-ray irradiation in Fig. 3 indicate that the borate group was the main structural unit in glass network. The infrared vibration at 992 cm⁻¹ corresponded to the B–O stretching vibration of tetrahedra BO₄ units in tri-, tetra- and pentaborate. While the B–O stretching of trigonal BO₃ units in tetrahedra BO₄ units were attributed to the vibration around 1122 cm⁻¹. The FTIR absorption around 1342 cm⁻¹ was assigned to the B–O stretching vibration of the trigonal (BO₃)₃⁻ units in meta-, pyro- and orthoborates. The vibration centered at 2923 and 2852 cm⁻¹ corresponded to the O–H stretching of hydroxyl OH⁻ groups, while the broad band around 3288 cm⁻¹ revealed to the vibration of OH⁻ groups and B–OH linkage. Before γ-irradiation, the vibration strength of these BO₄, BO₃, OH groups and B-OH linkage in 1.5CeGB glass were weaker than 0CeGB glass, indicating that the chemical groups in CeF₃-doped glasses have poorer connectivity compared to the undoped glass. After irradiation, the γ-rays could break some chemical bonds in the glass network, resulting to the decrease in vibration strength of those chemical complexes. The infrared absorption by such complexes then reduced which caused the increment of FTIR transmittance after irradiation. The change in vibration strength of 1.5CeGB glass due to γ-rays damage was less than that of the 0CeGB glass.
The optical properties. The transmission spectra of the Ce:GB glasses before γ-irradiation are shown in Fig. 4a. The unirradiated 0CeGB glass exhibited strong absorption in UV region with a transmission edge wavelength around 320 nm and the transmission spectra was shifted to longer wavelength with addition of CeF$_3$ content. This red-shift of glasses influenced by CeF$_3$ concentration were also found in several literatures\textsuperscript{18,34,35,38,46}. Considering on the effect of γ-rays, the 0CeGB glass after irradiation obviously absorbed photons in UV and VIS regions, as shown in Fig. 4b. The γ-rays generated the color center that increased the absorption in both regions, especially in the VIS range.

Figure 2. (a) The XANES spectra of Ce:GB glasses before and (b) after γ-irradiation, compared to the unirradiated CeF$_3$ and CeO$_2$ standard compounds.

Figure 3. The FTIR of 0CeGB and 1.5CeGB glass before/after γ-irradiation.
The increased absorption after irradiation is called "radiation-induced absorption" that can be considered from the change of transmittance before \( T_0 \) and after \( T \) irradiation at each wavelength by following relation 47, \[ T = T_0 \cdot e^{-\alpha_D x}, \] where \( \alpha_D \) is the radiation-induced absorption coefficient at each wavelength and \( x \) is the optical path length or thickness of sample. The \( \alpha_D \) were calculated in a range of VIS, 400–800 nm that the typical coefficient at 477 nm \( (\alpha_D-477) \) and the average coefficient \( (\alpha_D-\text{ave}) \) were shown in Table 2. The decrease of \( \alpha_D \) values represented the less effect of irradiation on glass with added CeF\(_3\) concentration, corresponding to the comparative transmission
as the amount of CeF₃ increased. Considering the effect of γ-rays, the refractive index of irradiated glasses were fluctuating with addition of CeF₃. To evaluate the total change of glass color that is damage from γ-rays, the color difference (\(\Delta E_{ab}\)) was calculated by equation 48:

\[
\Delta E_{ab} = \sqrt{(\Delta L')^2 + (\Delta a')^2 + (\Delta b')^2},
\]

where \(\Delta L', \Delta a'\) and \(\Delta b'\) is the difference values of such color parameters, before and after irradiation. The \(\Delta E_{ab}\) value decreased with increment of CeF₃ content corresponded to the change of glass color becoming more greenish yellow, as can be seen in Fig. 1. After irradiation, γ-rays decreased the \(L'^*\), changed the \(a'^*\) to be positive and increased the \(b'^*\), resulting to the dark tone color of undoped 0CeGB glass. For CeF₃-doped glasses after irradiation, the change of \(a'^*\) and \(b'^*\) tended to be less while the change of \(a'^*\) was fluctuating with addition of CeF₃. To evaluate the total change of glass color that is damage from γ-rays, the color difference (\(\Delta E_{ab}\)) between glasses before and after irradiation were calculated by equation 48.

\[
\Delta E_{ab} = \sqrt{(\Delta L'^*)^2 + (\Delta a'^*)^2 + (\Delta b'^*)^2},
\]

where \(\Delta L'^*, \Delta a'^*\) and \(\Delta b'^*\) is the difference values of such color parameters, before and after irradiation. The \(\Delta E_{ab}\) value decreased with increment of CeF₃ content corresponded to the change of glass color becoming more greenish yellow, as can be seen in Fig. 1. After irradiation, γ-rays decreased the \(L'^*\), changed the \(a'^*\) to be positive and increased the \(a'^*\), resulting to the dark tone color of undoped 0CeGB glass. For CeF₃-doped glasses after irradiation, the change of \(L'^*\) and \(b'^*\) tended to be less while the change of \(a'^*\) was fluctuating with addition of CeF₃. To evaluate the total change of glass color that is damage from γ-rays, the color difference (\(\Delta E_{ab}\)) between glasses before and after irradiation were calculated by equation 48.

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value as shown in Table 3. This means the electrons in glasses were more sensed and the molecules were easily polarized to such electric field by the increment of CeF₃ dopant and the γ-irradiation. The change of polarizability due to irradiation (Δαₘ) was smaller with adding CeF₃ concentration.

The αD-ave, ΔEab* and Δn value as a function of CeF₃ concentration were plotted cooperatively in Fig. 5 which those parameters owned the similar behavior on variation of CeF₃ content. This represents the ability of CeF₃ that enhanced the radiation hardness on glass optical properties.

The photoluminescence spectra and decay curves. The PL emission (solid line) and excitation (dash line) spectra under direct Ce³⁺ excitation of glasses before irradiation are shown in Fig. 6a. The luminescence intensity increased with increasing CeF₃ concentration in the range of 0.00–0.30 mol%. The intensity decreased for CeF₃ amounts larger than 0.30 mol% due to concentration quenching. The UV with 310 nm wavelength directly excited to the Ce³⁺ and promoted this ion from the ground 4f (2F⁵/₂) to the excited 5d state. The Ce³⁺ then relaxed to the lowest vibrational 5d state via a non-radiative relaxation (NR) process, followed by the 5d → 4f (2F⁵/₂) transition where a photon with 360 nm wavelength was emitted³⁴,³⁸,⁴⁰. After irradiation, the PL spectra in Fig. 6b shows a similar peak position as the spectra before irradiation. Concentration quenching was also observed when the CeF₃ doping concentration was more than 0.30 mol%. However, the luminescence intensity of the irradiated glasses decreased compared to the unirradiated ones. A clear evidence of intensity degradation is represented in the comparative spectra of 0.3CeGB glass in Fig. 6c. Moreover, an excitation peak around 275 nm of Gd³⁺ was also found and it overlapped on the left side of the Ce³⁺ excitation peak in those Fig. 6. The emission spectra of glasses under Gd³⁺ excitation were then studied and shown in Fig. 7a–c. The peak position, the influence of CeF₃ concentration and γ-irradiation on emission intensity were similar with the spectra under

Figure 5. The average radiation-induced absorption coefficient (αD-ave), color difference (ΔEab*) and refractive index difference (Δn) of Ce:GB glasses.

Figure 6. (a) The PL emission and excitation spectra under Ce³⁺ excitation of Ce:GB glasses before γ-irradiation and (b) after γ-irradiation, (c) the comparative PL spectra before and after γ-irradiation of 0.3CeGB glass.
direct Ce$^{3+}$ excitation. Additionally, a small peak of Gd$^{3+}$ emission under 6P7/2 → 8S7/2 transition was found at 312 nm wavelength. The strength of the Gd$^{3+}$ emission peak was weakened with increasing CeF$_3$ concentration because the excitation energy of Gd$^{3+}$ was more transferred to Ce$^{3+}$. The mechanism is as follows: the UV excitation with 275 nm excited the Gd$^{3+}$ from 8S7/2 to 6I7/2 state. NR then took Gd$^{3+}$ down to 6P7/2 level which was the intersection for next two separate routes. The first one was the 6P7/2 → 8S7/2 transition where Gd$^{3+}$ emitted the photon with 312 nm. For the second route, the energy transferred from 6P7/2 state of Gd$^{3+}$ to 5d state of Ce$^{3+}$. After that, the 5d → 4f (2F5/2) transition of Ce$^{3+}$ emitted the photon with 360 nm$^{34,38,40}$. Furthermore, there was a probability that the excitation with 275 nm also directly excited to Ce$^{3+}$ because its energy range of 5d state in glass is wide and overlaps with the 6I7/2 state of Gd$^{3+}$. This appeared as the overlapping between the excitation peak of Gd$^{3+}$ and Ce$^{3+}$ in the PL spectra. The possible mechanisms about the energy transition of Ce$^{3+}$ and Gd$^{3+}$ in the PL spectra are represented in Fig. 8.

The radioluminescence spectra and decay curves. The RL spectra before and after γ-irradiation in Fig. 10 show a strong emission from Ce$^{3+}$ around 381 nm wavelength. The incoming X-rays could initially inter-
act to the glass host. The X-rays energy then transferred to luminescence center (Gd$^{3+}$ or Ce$^{3+}$). After that, Gd$^{3+}$ could emit 312 nm luminescence but could not be detected in this experiment because of lower limit detection of spectrometer. There is also another possibility of Gd$^{3+}$ transferred the energy to Ce$^{3+}$ for luminescence under the 5d$\rightarrow$4f transition, similar process with the PL spectra. Additionally, the 5d–4f transition of Ce$^{3+}$ could be occurred from this scintillation process. The RL intensity of glasses tended to quench for CeF$_3$ concentrations that was greater than 0.30 mol%, like what was observed in the PL spectra. The γ-irradiation degraded the RL intensity of each glass which can be clearly observed in the comparative RL spectra of 0.3CeGB glass in Fig. 10c. The RL intensity of 0.3CeGB glass after irradiation decreased by 35% compared to its pre-irradiation intensity.

Discussion

Before γ-irradiation. In this part, only the influence of CeF$_3$ concentration on glasses before γ-irradiation are discussed and some explanations will be used continuously in the next subsection. The CeF$_3$ dopant possibly acted as a glass modifier that created non-bridging oxygen (NBO) and disrupted the connectivity of borate groups in the Gd$_2$O$_3$-B$_2$O$_3$ glass network. Consequently, the FTIR vibration strengths of those borate complexes in 1.5CeGB were weaker than in 0CeGB glass. The vibration of OH$^-$ groups in FTIR were also reduced by CeF$_3$ increment due to the reaction with F$^-$ ion as followed, 2OH$^-$ + 2F$^-$ → 2HF + in binary glass preparation as the glass O$_2$ $^{41,51}$. Since the transmission edge of host 0CeGB glass was overlapping to the 4f–5d transition of Ce$^{3+}$, the transmission spectra of CeF$_3$-doped glasses then were red-shifted by more influence of this transition via

Figure 9. (a) The decay curves under 290 nm excitation of Ce:GB glasses before γ-ray irradiation and (b) after γ-irradiation.

Figure 10. (a) The RL spectra of Ce:GB glasses before γ-irradiation $^{41}$ and (b) after γ-ray irradiation, (c) the comparative RL spectrum before and after γ-ray irradiation of 0.3CeGB glass.
increasing CeF₃ dopant. This red-shift caused the change of CIELAB parameters in Table 2 and the observed change of color becoming more greenish yellow in Fig. 1. Additionally, there were reported that the addition of NBO could increase the glass optical basicity and consequently affected to the red-shift in absorption–transmission spectra. The electrons at NBO sites in glass are less tightly bound and can be easily ionized by the electric field from incoming light, compared to electrons at bridging oxygen sites. The polarizability of glass then increased by following number of NBO with addition of CeF₃ content. Since the light was more sensed by the electrons, this light–electron interaction slowed down the speed of light (v) in glass resulting to the increment of refractive index by n = cv relation with added CeF₃ concentration.

The electronic configuration of Ce³⁺ ion is 4f⁰, which means that it has only one electron in the 4f shell to lose in order to have a more stable empty state. Therefore, Ce³⁺ can change to Ce⁴⁺ by losing one of its 4f electrons through the direct ionization, thermal ionization, or donation to hole by the process: Ce³⁺ + hole → Ce⁴⁺. On the other hand, a Ce⁴⁺ ion can accept an electron to form Ce³⁺ via the reaction: Ce⁴⁺ + electron → Ce³⁺[38,41]. This causes the coexistence of Ce³⁺ and Ce⁴⁺ ion in cerium doped materials, such as our glasses in this work. Since the electrons at NBO sites are less tightly bound, these electrons density could distribute and affect to the behavior of an 4f electron of Ce³⁺. The electrostatic pull between Ce³⁺ nucleus and its 4f orbital was weaken by such negative charges from NBO[41,55,56]. This increased the probability of Ce³⁺/Ce⁴⁺ ratio reduction with increasing of CeF₃ content as observed in the XANES spectra. Concentration quenching was found in both the PL and RL spectra of glasses doped with CeF₃ higher than 0.3 mol%. Quenching is due to the re-absorption of photons that are emitted by closely nearby Ce³⁺ neighbors. The shorter distance between Ce³⁺ ions and the dense ion distribution in glass provide this quenching effect, which also led to the reduction of decay time for glasses with more than 0.3 mol% of CeF₃.

After γ-irradiation. Considering the 0CeGB glass after irradiation, γ-rays could break some B–O–B and B–O–H linkages by following Eqs. (5) and (6), respectively.

\[ \text{B–O–B} \rightarrow \text{B–O–H} \] (5)

\[ \text{B–OH} \rightarrow \text{B} + \text{H}^⁻. \] (6)

Both “O–” and “O⋅” are the NBO, while the “B–” is the deformed borate complexes. Generally, the NBO and deformed borate are the charge defect which naturally pre-exist in the unirradiated metal-oxide borate glasses, also in our Gd₂O–B₂O₃ system, the γ-irradiation just increased the number of these complexes. The γ-rays could also ionize the chemical composition that generated the electron and hole in glass structure. This hole and electron could separate and move to trap with those charge defects in glass. Hole could be trapped by negative charge of NBO to form the non-bridging oxygen hole center (NBOHC). The electron was probably trapped by positive charge of deformed borate, becoming to the boron electron center (BEC). However, there was reported that the BEC in borate glass was unstable for temperature above 120 K and its number dramatically decreased to be negligible at about 320 K. Therefore, the main color center in our glasses after irradiation is NBOHC. This hole center is thought to absorb the photon around 3.8 eV (326 nm) and 2.6 eV (477 nm), that’s why the darkening and color change was obviously appeared in 0CeGB glass. This corresponded to the high value of radiation-induced absorption coefficient at 477 nm and its average value in VIS range of this glass as shown in Table 2. The radiation-induced absorption coefficients in UV range lower than 400 nm were not analyzed due to the overlapping of 4f–5d transition from Ce³⁺ on absorption of glass host. Since the irradiation possibly destroyed the B–O–B, B–O–H linkage and OH group shown in Eqs. (5) and (6), and disrupted the glass structure by formation of NBOHC, the FTIR vibration strength of 0CeGB glass then significantly decreased after irradiation. Moreover, the charge complexes such as the NBO and NBOHC created by γ-rays raised obviously the value of polarizability and refractive index in this glass.

For CeF₃-doped glasses after irradiation, the electron donation from Ce⁴⁺ to hole inhibited the hole trapping at the charge defect site such as NBO. Consequently, the number of NBOHC was decreased and the structure of 1.5CeGB glass was more conserved from the disruption than 0CeGB glass, as shown in FTIR spectra. The NBOHC reduction with the addition of CeF₃ content also caused a decrease of those radiation damage parameters such as the radiation-induced absorption coefficient, the color difference, the change of polarizability and refractive index. Especially in 1.5CeGB glass, these values were close to zero which represented the highest radiation hardness.

The donation of an electron from Ce⁴⁺ to hole caused the reduction of Ce³⁺/Ce⁴⁺ ratio in 0.05CeGB and 0.3CeGB glass after irradiation, observed by XANES spectra. For 1.5CeGB glass, the large amount of CeF₃ dopant created high number of pre-existing NBO in the glass network, and there were the electrons created by irradiation those could not trap to BEC because this center was unstable as previously mentioned. Some Ce⁴⁺ possibly accepted an electron from NBO and unstable BEC which changed this ion back to Ce³⁺. The Ce³⁺/Ce⁴⁺ ratio in 1.5CeGB glass therefore slightly increased by irradiation. The PL and RL luminescence intensity of CeF₃-doped glasses decreased after irradiation due to the reduction of Ce³⁺/Ce⁴⁺ ratio. Since the absorption energy of defect (NBOHC at 3.8 eV) in UV region overlapped to the 4f–5d transition of Ce³⁺ in this glass, the UV excitation energy on decay time measurement was possibly trapped by defect, resulting to longer decay time after irradiation.

From all results, the 0.3CeGB glass is a promising new glass scintillator, with the highest emission intensity among the glasses studied in this work, a relatively fast nanoseconds decay time and excellent radiation hardness.
Conclusion
Various properties of CeF$_3$-doped Gd$_2$O$_3$-B$_2$O$_3$ glasses before and after γ-irradiation were comparatively investigated. XANES results show that the major and minor oxidation states of cerium ion in glasses were Ce$^{4+}$ and Ce$^{3+}$, respectively. Before irradiation, the analysis of glasses' transparency, FTIR, refractive index and polarizability indicated that CeF$_3$ degraded the connectivity and possibly created NBO in glass structure. This NBO caused the reduction of Ce$^{3+}$/Ce$^{4+}$ ratio, the red-shift in transmission spectra and the raise of refractive index with addition of CeF$_3$ content. That red-shift also was influenced by 4f–5d transition of Ce$^{3+}$ dopant. After irradiation, γ-rays damaged the glass structure, broke the chemical bond, and created the color center in the borate network former. That center is NBOHC which absorbed photons in VIS region, resulting to the darkening and color change in glasses after irradiation. Moreover, the polarizability and refractive index of glasses were increased by the formation of NBO and NBOHC generated by irradiation. The addition of CeF$_3$ concentration in glass relieved these irradiation effects. Due to the electron donation from Ce$^{3+}$ to hole, number of NBOHC were annihilated. The radiation damage indicators such as the radiation-induced absorption coefficient, the color difference, the change of polarizability and refractive index then decreased in value with increasing CeF$_3$ dopant. These results confirm the ability of CeF$_3$ that enhances the radiation hardness of glass. The PL and RL intensity of CeF$_3$-doped glasses decreased after irradiation due to the reduction of Ce$^{3+}$/Ce$^{4+}$ ratio via electron donation of Ce$^{3+}$. The decay times of glasses after irradiation were longer, compared to them before irradiation because the excitation energy was possibly trapped by defect (NBOHC). The Gd$_2$O$_3$-B$_2$O$_3$ glass doped with 0.30 mol% of CeF$_3$ exhibited the highest emission intensity, fast 24–27 ns decay time and owned the radiation hardness property, making it a promising new glass scintillator.

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References
1. Yanagida, T. et al. Study of the correlation of scintillation decay and emission wavelength. Radiat. Meas. 55, 99–102 (2013).
2. Yanagida, T. Inorganic scintillating materials and scintillation detectors. Proc. Jpn. Acad. Ser. B Phys. Biol. Sci. 94, 75–97 (2018).
3. Arikawa, Y. et al. Pr$^{3+}$-doped fluorine oxide lithium glass as scintillator for nuclear fusion diagnostics. Rev. Sci. Instrum. 80, 113504 (2009).
4. Watanabe, K. et al. Pr or Ce-doped, fast-response and low-afterglow cross-section-enhanced scintillator with 6Li for down-scattered neutron originated from laser fusion. J. Cryst. Growth 362, 288–290 (2013).
5. Yamanoye, K. et al. Luminescence properties of Nd$^{3+}$ and Er$^{3+}$ doped glasses in the VUV region. Opt. Mater. (Amst.) 35, 1962–1964 (2013).
6. Oyarzun, D. F. et al. The effects of CeF$_3$ on polarization and optical, transmission and photon/neutron attenuation properties of boron-tellurite glasses. J. Non. Cryst. Solids 544, 120171 (2020).
7. Al-Buriahi, M. S. et al. Effect of CeF$_3$ doping on the radiation attenuation properties of tellurite glasses containing V$_2$O$_5$ and Nb$_2$O$_5$. Appl. Phys. A Mater. Sci. Process. 127, 1–12 (2021).
8. Hegazy, H. H. et al. The effects of CeF$_3$ on polarization, optical transmission, and photon/neutron attenuation properties of boron-tellurite glasses. J. Inorg. Organomet. Polym. Mater. 31, 2331–2338 (2021).
9. Al-Buriahi, M. S., Somaily, H. H., Al-Buriahi, M. S. Effect of CeF$_3$-doped Ag$_2$O/N$_2$O substitution on the radiation shielding ability of tellurite glass system via XCOM approach and FLUKA simulations. Phys. Scr. 96, 065308 (2021).
10. Al-Buriahi, M. S., Al-Buriahi, H. T., Alrowaili, Z. A. Optical and radiation shielding studies on tellurite glass system containing ZnO and NiO. Ceram. Int. 46, 19078–19083 (2020).
11. Al-Buriahi, M. S., Alrowaili, Z. A. Mechanical and radiation shielding properties of tellurite glasses doped with ZnO and NiO. Ceram. Int. 46, 19078–19083 (2020).
12. Al-Buriahi, M. S., Singh, V. P., Alalawi, A., Sridharan, C. & Tonguc, B. T. Optical and radiation shielding properties of TeO$_2$–Ag$_2$O-WO$_3$ glasses. Ceram. Int. 156, 154–15472 (2020).
13. Al-Buriahi, M. S. et al. Effect of chromium oxide on the physical, optical, and radiation shielding properties of lead sodium borate glasses. J. Non. Cryst. Solids 544, 120171 (2020).
14. Al-Buriahi, M. S. et al. Effect of Nb$_2$O$_5$ and V$_2$O$_5$ on the radiation attenuation properties of tellurite glasses containing V$_2$O$_5$ and Nb$_2$O$_5$. Appl. Phys. A Mater. Sci. Process. 127, 1–12 (2021).
15. Al-Buriahi, M. S. et al. Optical and radioluminescent behaviors of Sm$^{3+}$ doped high-density tungsten gadolinium borate scintillating glass. J. Alloy. Compd. 849, 108350 (2019).
16. Van Eijk, C. W. E. Inorganic scintillators for thermal neutron detection. Radiat. Meas. 38, 337–342 (2004).
17. Lee, G. H., Chang, Y. & Kim, T.-J. Properties and Possible Application Areas. Ultrasound Lanthanide Oxide Nanoparticles for Biomedical Imaging and Therapy (Woodhead Publishing, 2014).
18. Xin, P. et al. Study on the sensitization of Gd$^{3+}$ on Ce$^{3+}$/Tb$^{3+}$ co-doped GBS scintillating glass. J. Non. Cryst. Solids 501, 411–416 (2020).
19. Lai, Y. et al. Investigation of gammray induced optical property changes in non-doped and Ce-doped lithium-rich oxide glass. Radiat. Phys. Chem. 179, 109272 (2021).
20. Wantana, N. et al. Investigation of gamma-ray induced optical property changes in non-doped and Ce-doped lithium-rich oxide glass. Radiat. Phys. Chem. 179, 109272 (2021).
21. Wantana, N. et al. X-ray/proton and photoluminescence behaviors of Sm$^{3+}$ doped high-density tungsten gadolinium borate scintillating glass. J. Alloy. Compd. 849, 156574 (2020).
22. Van Eijk, C. W. E. Inorganic scintillators for thermal neutron detection. Radiat. Meas. 38, 337–342 (2004).
23. Lee, G. H., Chang, Y. & Kim, T.-J. Properties and Possible Application Areas. Ultrasound Lanthanide Oxide Nanoparticles for Biomedical Imaging and Therapy (Woodhead Publishing, 2014).
24. Xu, P. et al. Study on the sensitization of Gd$^{3+}$ on Ce$^{3+}$/Tb$^{3+}$ co-doped GBS scintillating glass. J. Non. Cryst. Solids 501, 411–416 (2018).
25. Wantana, N. et al. Investigation of gamma-ray induced optical property changes in non-doped and Ce-doped lithium-rich oxide glass. Radiat. Phys. Chem. 179, 109272 (2021).
26. Som, T. & Karmakar, B. Nephelauxetic effect of low phonon antimony oxide glass in absorption and photoluminescence of rare-earth ions. Spectrochim. Acta A Mol. Biomol. Spectrosc. 79, 1766–1782 (2011).
27. Wantana, N. et al. Tunable orange, yellow and white emission of Pr$^{3+}$-doped tungsten gadolinium borate glasses. J. Non. Cryst. Solids 554, 120603 (2021).
28. Taki, Y., Shinohara, K., Homma, T., Dimitrov, V. & Komatsu, T. Electronic polarization and interaction parameter of gadolinium tungsten borate glasses with high WO3 content. J. Solid State Chem. 220, 191–197 (2014).
29. Wantana, N. et al. High density tungsten gadolinium borate glasses doped with Eu$^{3+}$ ion for photonic and scintillator applications. Radiat. Phys. Chem. 172, 108868 (2020).
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**Author contributions**
E.K. analyzed the results, made the discussion, and edited the manuscript. N.W. measured and categorized the glasses properties. Y.R. prepared the glasses and found the optimum condition for synthesis. M.C.-R. and K.Y. provided the γ-irradiation on glasses and discussed its influence on the glasses properties. P.K. studied and analyzed the XANES spectra of glasses. H.J.K. and J.K. designed and gave the direction of research, including the consultation to develop this work.

**Competing interests**
The authors declare no competing interests.

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