Observed CN and HCN intensity ratios exhibit subtle variations in extreme galaxy environments

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ABSTRACT

We use both new and archival ALMA data of three energy lines each of CN and HCN to explore intensity ratios in dense gas in NGC 3256, NGC 7469, and IRAS 13120-5453. The HCN (3-2)/HCN (1-0) intensity ratio varies in NGC 3256 and NGC 7469, with superlinear trends of 1.53 ± 0.07 and 1.55 ± 0.05, respectively. We find an offset to higher HCN (3-2)/HCN (1-0) intensity ratios (~0.8) in IRAS 13120-5453 compared to NGC 3256 (~0.3 – 0.4) and NGC 7469 (~0.3 – 0.5). The HCN (4-3)/HCN (3-2) intensity ratio in NGC 7469 has a slope of 1.34 ± 0.05. We attribute the variation within NGC 3256 to excitation associated with the northern and southern nuclei. In NGC 7469, the variations are localized to the region surrounding the active galactic nucleus. At our resolution (~700 pc), IRAS 13120-5453 shows little variation in the HCN intensity ratios. Individual galaxies show nearly constant CN (2-1)/CN (1-0) intensity ratios. We find an offset to lower CN (2-1)/CN (1-0) intensity ratios (~0.5) in NGC 3256 compared to the other two galaxies (~0.8). For the CN (3-2)/CN (2-1) intensity ratio, NGC 7469 has a superlinear trend of 1.55 ± 0.04, with the peak localized toward the active galactic nucleus. We find high (~1.7) CN (1-0)/HCN (1-0) intensity ratios in IRAS 13120-5453 and in the northern nucleus of NGC 3256, compared to a more constant ratio (~1.1) in NGC 7469 and non-starbursting regions of NGC 3256.

Key words: ISM: molecules – galaxies: ISM – galaxies: starburst – galaxies: nuclei

1 INTRODUCTION

In dense regions of molecular clouds, observers use molecules with high critical densities such as HCN, HCO+, and CS to identify regions of active and future star formation (Wu et al. 2010; Kennicutt & Evans 2012). In particular, HCN luminosity has been shown to correlate well with total far-infrared (FIR) luminosity, a tracer of the star formation rate (SFR) of star-forming galaxies (SFR = M_{IR}/τ, Vollmer et al. 2017; Privon et al. 2017).

An interesting molecule that is astrochemically related to HCN is the cyanide radical (CN). CN also has a high critical density and is primarily formed from photodissociation of HCN and neutral-neutral reactions with N, C2, CH2, and CH (Aalto et al. 2002; Boger & Sternberg 2005; Chapillon et al. 2012). Intermediate stages in the reaction pathways involve neutral and ionized carbon (C and C+; Boger & Sternberg 2005). CN is thus thought to preferentially form in regions illuminated by intense radiation fields, including the ultra-violet (UV) radiation fields of photo-dissociation regions (PDRs) surrounding massive stars (Fuente et al. 1993, 1995; Greaves & Church 1996; Bachiller et al. 1997; Rodriguez-Franco et al. 1998; Boger & Sternberg 2005; Ginard et al. 2015). CN will also be more likely to form when exposed to radiation in X-ray dominated regions (XDRs) near AGN (Meijerink & Spaans 2005; Meijerink et al. 2007; García-Burillo et al. 2010) and increased cosmic ray ionization rates (Boger & Sternberg 2005; Bayet et al. 2011; Aladro et al. 2013). The CN/HCN abundance ratio is therefore predicted to increase where photodestruction rates of HCN are high and carbon is ionized, and to decrease with cloud depth as HCN is protected from external radiation (Sternberg & Dalgarno 1995; Rodriguez-Franco et al. 1998; Boger & Sternberg 2005).

Galaxies with starbursts and/or AGN are prime environments for probing CN and HCN and their interactions with radiation fields. U/LIRGs are typically interacting or merging galaxies characterized by high infrared luminosities (Sanders et al. 2003), large fractions of dense molecular gas (Solomon et al. 1992; Gao & Solomon 2004; Privon et al. 2017; Sliwa et al. 2017), and high star formation rate surface densities (Σ_{SFR}, Vollmer et al. 2017; Privon et al. 2017). Additionally, many of these systems host AGN, indicating the presence of high luminosity black holes and intense UV radiation fields. The CN/HCN ratio may therefore provide insights into the complex interplay between star formation and AGN activity in these systems.

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Table 1. Basic properties of the galaxy sample.

| Propertya | NGC 3256 | NGC 7469 | IRAS 13120 |
|-----------|----------|----------|------------|
| Luminosity class | LIRG | LIRG | U/LIRG |
| log($L_{\text{IR}}$) ($L_{\odot}$) | 11.75 | 11.60 | 12.29 |
| AGN contribution | < 5% | 32 – 40% | 17 – 33% |
| RA (J2000) | 10h27m51.3s | 23h03m15.6s | 13h15m06.4s |
| DEC (J2000) | +43°54′13.5″ | +08°52′26″ | −55°09′22.6″ |
| Diameter | 3.8′ × 2.1′ | 1.5′ × 1.1′ | 0.33′ × 0.33′ |
| Redshift | 0.00935 | 0.01632 | 0.03076 |
| $D_1$ (Mpc)b | 44 | 66 | 134 |
| $\langle SFR \rangle$ (M$_\odot$ yr$^{-1}$)c | 84 | 60 | 292 |

Notes: 
a) NGC 3256 and NGC 7469 properties are retrieved from the NASA/IPAC Extragalactic Database (NED). IRAS 13120 properties come from Simbadd).
b) $L_{\text{IR}}$ data from Sanders et al. (2003) and are corrected for luminosity distance.

c) The AGN contribution to $L_{\text{IR}}$ is estimated from 6 µm and 24 µm emission for NGC 3256 and NGC 7469 (Alonso-Herrero et al. 2012). For IRAS 13120, the AGN contribution is estimated from 15 µm and 30 µm (Veilleux et al. 2013), 60 µm (Teng et al., 2015), and total 8-1000 µm emission (Iwasawa et al. 2011).
d) Luminosity distances from redshifts (corrected to the 3K CMB reference frame) and assuming H$_0$ = 70.5 km s$^{-1}$ Mpc$^{-1}$. For NGC 7469, the SN Type Ia distance is from Ganesalingam et al. (2013).
e) SFRs calculated from $L_{\text{IR}}$ using Equation 12 in Kennicutt & Evans (2012); not corrected for AGN contribution.

2 OBSERVATIONS AND DATA REDUCTION

2.1 U/LIRG sample

Our sample consists of two LIRGs, NGC 3256 and NGC 7469, and one U/LIRG, IRAS 13120-5453 (hereafter referred to as IRAS 13120). The individual properties of these galaxies are summarized in Table 1. The different galaxy environments available in this sample allow us to probe the effects of both UV and X-ray radiation fields on the molecular gas.

NGC 3256 is a nearby luminous galaxy merger that can be separated into northern and southern nuclear regions. The northern nucleus is clearly face-on and has a nuclear disk with significant starburst activity (Sakamoto et al. 2014). The southern nucleus is mostly edge-on and thought to host an embedded dormant AGN (Sakamoto et al. 2014). NGC 3256 is a complicated system containing molecular outflows from both nuclei (Sakamoto et al. 2014) and high fractions of shocked and dense gas (Harada et al. 2018). Sakamoto et al. (2014) find that the outflows in the southern nucleus are highly collimated, bipolar nuclear jets with velocities of ~2000 km s$^{-1}$. Brunetti et al. (2021) argue that, due to a lack of significant trends in molecular gas surface density, brightness temperature, and velocity dispersion with physical scale, NGC 3256 must contain a smooth interstellar medium down to ~55 pc scales.

NGC 7469 hosts a nuclear luminous type-1 AGN ($L_{2-10keV}$ = 1.5$x_{10^{43}}$ erg s$^{-1}$; Liu et al. 2014) and has a starburst ring surrounding the central portion of the galaxy (Izumi et al. 2015, 2020). The AGN is creating an XDR ~50 pc in size in the nuclear region, and this XDR is surrounded by a circumnuclear disk of star-forming gas (Izumi et al. 2020). Izumi et al. (2015) find the integrated intensities of HCN (4-3), HCO$^+$ (4-3), and CS (7-6) are higher in the nuclear than in the starburst ring and predict that the HCN abundance is enhanced in the nucleus due to increased sub-mm radiation near the AGN (Izumi et al. 2013, 2016).

IRAS 13120 has been optically classified as a Seyfert 2 galaxy (Véron-Cetty & Véron 2010). Teng et al. (2015) find that observations of IRAS 13120 are consistent with the galaxy hosting an inactive, Compton-thick AGN ($N_H > 10^{24}$ cm$^{-2}$). 12CO observations by Sliwa et al. (2017) show evidence for a young starburst ($<$ 7 Myr) in the central ~500 pc region. Privon et al. (2017) find the HCN/HCO$^+$ ratio is higher in the central starburst and suggest this is from mechanical heating of the gas by supernovae feedback.

2.2 Data and imaging

Table 2 lists the project IDs for the new and archival ALMA data used in our analysis. The data were reduced and imaged with the Common Astronomy Software Application (CASA; McMullin et al. 2007). The raw $uv$ data were calibrated for each project using the relevant CASA version. All subsequent data reduction was performed using CASA version 5.6.1. Continuum subtraction was performed on each $uv$ dataset using line-free channels and CASA’s uvcontsub task.

Given the large collection of ALMA observations from which our data were obtained, we put significant effort into matching the spatial and spectral resolutions of the lines. The $uv$ data for each galaxy were imaged individually with CASA’s tclean task. All imaging was done using Briggs’ weighting (Briggs 1995). Making use of tclean’s cell, uv taper, uv range, and width parameters, the $uv$ coverage, spectral channel width, and dirty beam sizes were compared and matched for each galaxy. The limiting factors were the largest minimum coverage, the largest spectral channel width, and the largest dirty beam size. We note that applying a $uv$ cut-off and...
Table 2. New and archival ALMA projects with HCN (J = 1-0, 3-2, 4-3) and CN (N = 1-0, 2-1, 3-2) observations of the three target galaxies.

| Line          | Project code         | PI       | ALMA data reference   |
|---------------|----------------------|----------|-----------------------|
| NGC 3256      | 2015.1.00993.S       | Michiyama, T. | Michiyama et al. (2018) |
| HCN (1-0)     | 2015.1.00412.S       | Harada, N.  | Harada et al. (2018)  |
| HCN (4-3)     | 2018.1.00493.S       | Wilson, C.  | This paper             |
| CN (1-0)      | 2011.0.00525.S       | Sakamoto, K. | Sakamoto et al. (2014) |
| CN (2-1)      | 2015.1.00412.S       | Harada, N.  | Harada et al. (2018)  |
| CN (3-2)      | 2018.1.00493.S       | Wilson, C.  | This paper             |
| NGC 7469      | 2012.1.00165.S       | Izumi T.   | Izumi et al. (2015)   |
| HCN (1-0)     | 2012.1.00034.S       | Himanish M. | Himanish et al. (2016) |
| HCN (4-3)     | 2013.1.00218.S       | Izumi T.   | Izumi et al. (2015)   |
| CN (1-0)      | 2015.1.00412.S       | Harada, N.  | Harada et al. (2018)  |
| CN (3-2)      | 2018.1.00493.S       | Wilson, C.  | This paper             |
| IRAS 13120    | 2013.1.00379.S       | Sliwa, K.  | Sliwa et al. (2017)   |
| HCN (1-0)     | 2018.1.00493.S       | Wilson, C.  | This paper             |
| HCN (4-3)     | 2015.1.00102.S       | Iono D.    | Fluetisch et al. (2019) |
| CN (1-0)      | 2016.1.00287.S       | Sliwa, K.  | Wilson et al. (2019)  |
| CN (2-1)      | 2018.1.00777.S       | Sliwa, K.  | –                     |
| CN (3-2)      | 2018.1.00493.S       | Wilson, C.  | This paper             |

Notes: aReference uses the CO (1-0) transition from these project IDs.
bReference uses the CO (3-2) transition from this project ID.

taper with the uvr_{range} and uvtaper parameters in tclean will limit the number of short baselines in our data and thus could potentially lead to missing flux on the largest angular scales. The uvr_{range} parameter was matched between all lines in each galaxy in order to recover similar flux scales for each line. We tried to limit the problem of missing flux in our interferometric observations by choosing the largest minimum uvr_{range} cut-off while still matching all lines.

A measure of the mean RMS noise was found from both the dirty and final smoothed image cubes for each spectral line in CASA using the imstat task on line-free channels. The sensitivities varied for each line in each galaxy depending on the specific ALMA observations and can be found in Tables 3 and A1. The imaging of all galaxies was limited to regions where the primary beam response was greater than 20% and cleaning was performed using CASA’s automatic-thresholding (Kepley 2019) down to the 2σ level using the line sensitivities from the dirty cubes. Masks produced by tclean’s auto-masking algorithm were checked for consistency during major cleaning cycles. Most observations were completed with a single pointing and imaged with matched phase centres in tclean, except for the ALMA Band 7 observations of CN (3-2) and HCN (4-3) in NGC 3256 which had multiple pointings. In these cases, the imaging was done as a mosaic with the same matched phase centres. A final smoothing to a common round beam was completed for all lines in each galaxy with the CASA task imsmooth.

2.2.1 NGC 3256 processing

The spectral resolution in NGC 3256 was limited by the CN (1-0) line at 26.43 km s\(^{-1}\) and all lines were binned in velocity to this channel width. The limiting dirty beam size was set by the HCN (1-0) observations at 2.1\(^{\prime\prime}\) × 1.7\(^{\prime\prime}\), PA = 88°. The parameter uvtaper was used to taper the longest baselines in each dataset until the beams were matched relatively well, before a final smoothing to a 2.2\(^{\prime\prime}\) round beam was applied. For the HCN (3-2), HCN (4-3), CN (2-1), and CN (3-2) lines, natural weighting (Briggs’s with robust = 2.0) was used to increase sensitivity and naturally increase the dirty beam size before tapering. For HCN (1-0) and CN (1-0), the resolution was close enough to the target size of 2.2\(^{\prime\prime}\) that little tapering was required and Briggs’s weighting with robust = 0.5 was sufficient.

2.2.2 NGC 7469 processing

The three CN transitions in NGC 7469 had limiting spectral resolutions of \~1.57 km s\(^{-1}\); however, the channel widths were increased to an integer multiple of 4 \times 1.57 = 20.68 km s\(^{-1}\) to smooth the data spectrally and increase the signal-to-noise. This smoothing helped match the spectral resolution of the NGC 3256 and IRAS 13120 data. The CN (1-0) line limited our resolution to a beam size of 0.89\(^{\prime\prime}\) × 0.55\(^{\prime\prime}\), PA = -47.3°. We thus targeted a smoothed 0.95\(^{\prime\prime}\) round beam, and adjusted weighting and uvtaper parameters to match this target. HCN (1-0), CN (1-0), and CN (3-2) were imaged using robust = 0.5, HCN (3-2), HCN (4-3), and CN (3-2) were imaged using natural weighting.

2.2.3 IRAS 13120-5453 processing

In the IRAS 13120 data, the native spectral resolution varied for each of the three CN lines. The HCN lines all had the same \~3.3 km s\(^{-1}\) velocity resolution. Channel widths were fixed at 20.64 km s\(^{-1}\), an integer multiple of the CN (1-0) transition line. This spectral smoothing enhanced the signal-to-noise ratio for the large velocity dispersions seen in IRAS 13120. The limiting beam size was the HCN (1-0) line at 1.01\(^{\prime\prime}\) × 0.51\(^{\prime\prime}\), PA = 73.7°, imaged with robust = 0.5. The remaining lines were imaged with robust = 2.0, as they all had dirty beams with better than 0.6\(^{\prime\prime}\) resolution. Tapers were applied and we smoothed the data to a final target 1.1\(^{\prime\prime}\) round beam.

Table 3. Data reduction imaging properties and line sensitivities.

| Imaging property | NGC 3256 | NGC 7469 | IRAS 13120 |
|------------------|----------|----------|------------|
| Beam size\(^a\)  | 2.2\(^{\prime\prime}\) | 0.95\(^{\prime\prime}\) | 1.1\(^{\prime\prime}\) |
| Beam (pc)        | 469      | 304      | 715        |
| Velocity resolution (km s\(^{-1}\))\(^b\) | 26.43    | 20.68    | 20.64      |
| Pixel size\(^b\) | 0.3\(^{\prime\prime}\) | 0.15\(^{\prime\prime}\) | 0.15\(^{\prime\prime}\) |
| Re-binned pixel size\(^b\) | 1.1\(^{\prime\prime}\) | 0.475\(^{\prime\prime}\) | 0.55\(^{\prime\prime}\) |
| uvr_{range} cut-off (kλ)\(^c\) | > 15    | > 19.8   | > 15.9     |
| Maximum recoverable scale\(^c\) | 8.25\(^{\prime\prime}\) | 6.25\(^{\prime\prime}\) | 7.78\(^{\prime\prime}\) |

Notes: aBeams smoothed, rounded, and matched to this resolution for all lines.
bThe velocity resolution and pixel size were matched between all lines.
cuvr_{range} cut-off based on the minimum uv range covered by all lines in each galaxy and has also been converted to a maximum recoverable scale of emission.

dSensitivities determined using line-free channels of dirty image cubes.
Table 4. CN and HCN measured intensity ratios.

| Region   | HCN (3-2) | HCN (4-3) | CN (2-1) | CN (3-2) | CN (1-0) | CN (3-2) |
|----------|-----------|-----------|----------|----------|----------|----------|
|          | HCN (T=0) | HCN (T=2) | CN (T=0) | CN (T=2) | CN (T=0) | CN (T=2) |
| N. nucl. | 0.39(4)   | 0.48(7)   | 0.47(5)  | 0.27(4)  | 1.7(1)   | 0.55(8)  |
| S. nucl. | 0.33(4)   | 0.46(7)   | 0.44(5)  | 0.22(3)  | 1.12(9)  | 0.34(5)  |
| Non-nucl.| 0.27(3)   | 0.45(6)   | 0.43(5)  | 0.21(3)  | 1.14(8)  | 0.38(5)  |

NGC 7469

| Nucl.    | 0.52(6)   | 0.48(7)   | 0.80(9)  | 0.45(6)  | 1.13(9)  | 0.8(1)   |
| Non-nucl.| 0.33(4)   | 0.37(5)   | 0.76(9)  | 0.30(4)  | 1.11(8)  | 0.8(1)   |

IRAS 13120

| Global   | 0.80(9)   | 0.8(1)    | 0.77(9)  | 0.53(8)  | 1.7(1)   | 0.8(1)   |
| Peak     | 0.79(9)   | 0.9(1)    | 0.80(9)  | 0.57(8)  | 1.5(1)   | 0.8(1)   |

Notes: Uncertainties are measurement plus calibration uncertainties and are given as the uncertainty on the last digit (i.e. 0.39(4) = 0.39 ± 0.04). ALMA calibration uncertainties are 5% (Band 3) and 10% (Band 6 and 7).

2.3 Integrated intensities and measured ratios

We produced integrated intensity (moment 0) maps of all lines in our three galaxies. Cleaned, smoothed data cubes were extracted from CASA and subsequent processing, including the creation of the moment 0 maps, was done using the astropy3 software in Python 3.7. The cubes were trimmed to include only channels with line emission (found using the CASA viewer). The HCN (1-0), HCN (3-2), and HCN (4-3) line cubes were trimmed to include the same range in velocity. The CN lines were trimmed depending on their hyperfine structure so that all hyperfine lines were included in the calculation of the integrated intensities. Exact velocity ranges for all six lines can be found in Table A1. An RMS cut-off of 3.5σ was used to mask the cleaned cubes to include strong emission while avoiding noisy emission in lines with lower S/N. This 3.5σ limit used the line sensitivities calculated in the smoothed cubes and was chosen to overcome the varying sensitivities between different lines in each galaxy. The moment 0 maps were further masked to only include pixels with emission in all six spectral lines. In cases where multiple hyperfine lines are present (the HCN (1-0), CN (1-0), CN (2-1), and CN (3-2) lines), the integrated intensities include all hyperfine lines that were observed. The moment 0 maps were converted to physical units of K km s⁻¹ and corrected for the primary beam response. Moment 0 maps for each line can be found in the Appendix (Figures A1, A2, and A3). Uncertainty maps were made to match each moment 0 map. Individual uncertainties on each pixel were calculated as

\[ \sigma = \text{rms} \cdot \sqrt{\frac{\Delta V_{\text{chan}}}{V_{\text{line}}}} \]

where \( \Delta V_{\text{chan}} \) is the velocity width of an individual channel and \( V_{\text{line}} \) is the width of the line integrated in that pixel. The integrated intensity maps were used to calculate total intensities and intensity ratios. Maps of the ratios covering the high S/N regions in each galaxy are shown in Figures 1 and 2.

We calculate global line ratios for each galaxy using the average line intensities in K km s⁻¹. Along with the global ratios, we measured line ratios in the northern and southern nuclei of NGC 3256 and the nucleus of NGC 7469. The central nuclear pixels were identified using the 93 GHz radio continuum emission peak for NGC 3256 and NGC 7469. To isolate the nuclear regions, we used an aperture centred on the nuclear pixel. The size of the aperture is equal to the full-width at half-maximum of the beam in each galaxy. Additionally, for NGC 3256 and NGC 7469 we calculated non-nuclear line ratios using all non-nuclear pixels. For IRAS 13120, we measured the line ratios at a single pixel at the 93 GHz radio continuum peak. Table 4 presents the results of this analysis and the measured ratios in the different regions of each galaxy.

2.4 Pixel-by-pixel comparisons

We explored correlations between spectral lines using a pixel-by-pixel comparison of the intensities. We first re-sampled the imaged data cubes to half the full-width at half-maximum of the beam. The number of pixels sampled across the beam was reduced by a factor of 3 for all galaxies, and the new pixel sizes can be found in Table 3. Following the same process as described in Section 2.3, moment 0 maps with associated uncertainties were used to make masked, matched pixels in the re-sampled image cubes in physical units of K km s⁻¹. For NGC 7469, we corrected the measured intensities in the scatter plots for an inclination angle of 45 ± 5 degrees (Davies et al. 2004). The intensities of the other two galaxies were not corrected for inclination angle⁴.

We present pixel-by-pixel comparisons of the integrated intensities in the scatter plots of Figures 3 and 4. We used the Limix linear regression method to fit the integrated intensity pixels in log-log space and account for the higher uncertainties seen in the lower S/N pixels (Kelly 2007; Meyers 2018). The resulting slopes found for each intensity ratio are presented in Table 5. Full details of all calculated slopes and intercepts for all line ratio combinations can be found in Table A2 in the Appendix.

3 VARIATIONS IN THE CN AND HCN INTENSITY RATIOS

In this section, we present the measured intensity ratios and slopes from our CN and HCN line analysis. Discussion and physical interpretation of the ratios is deferred until Section 4.

4 NGC 3256 is a complicated system where the northern and southern nuclei are fairly face-on and edge-on, respectively (i(ν) ≈ 30° and i(ν) ≈ 80°; Sakamoto et al. 2014), making it challenging to correct for any specific angle. No inclination angle for IRAS 13120 was found in the literature. We performed the same analysis applying a conservative estimate for inclination angles of 60° to NGC 3256 and IRAS 1320 and found that the resulting slopes and intercepts match the results in Table 5 within one standard deviation.

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3.1 Regions of enhanced HCN ratios

We find variations in the HCN (3-2)/HCN (1-0) ratio in NGC 3256 and NGC 7469 (Figure 3a), with spatial variations in the ratio within these galaxies (Figures 1a and 1e). Both galaxies show a similar superlinear trend (1.53 ± 0.07 for NGC 3256, 1.55 ± 0.05 for NGC 7469). The northern nucleus of NGC 3256 shows a marginal increase in the HCN (3-2)/HCN (1-0) ratio to a value of 0.39 ± 0.04, in contrast to the value of 0.27 ± 0.03 in the non-nuclear pixels. In NGC 7469, the HCN (3-2)/HCN (1-0) ratio is significantly higher (0.52 ± 0.06) in the central ~ 500 pc compared to the value of the non-nuclear pixels (0.33 ± 0.04) in the rest of the disk.

Comparing the two galaxies, the HCN (3-2)/HCN (1-0) and HCN (4-3)/HCN (1-0) global and non-nuclear ratios in NGC 3256 and NGC 7469 are similar within our uncertainties (Table 4 and A3). In the nuclear regions, however, both ratios are higher in NGC 7469 than in NGC 3256.

There is no difference in the HCN (3-2)/HCN (1-0) ratio for IRAS 13120 between the global ratio and the continuum peak. IRAS 13120 shows a slight sublinear trend in the HCN (3-2)/HCN (1-0) ratio, with a slope of 0.84 ± 0.06 (Figure 3a), significantly different from the superlinear trends with slopes > 1.5 found in NGC 3256 and NGC 7469. The HCN (3-2)/HCN (1-0) ratio of the ULIRG is higher (~0.8) than both the global and nuclear values of the LIRGs (~0.3 – 0.5).

The HCN (4-3)/HCN (3-2) ratio in NGC 3256 has a slope consistent with unity (0.99 ± 0.04). This ratio in NGC 3256 has nearly identical values of 0.48 ± 0.07, 0.46 ± 0.07, and 0.45 ± 0.06 in the northern nucleus, southern nucleus, and non-nuclear pixels, respectively. Slight variations can be seen in the HCN (4-3)/HCN (3-2) ratio in NGC 3256 in Figure 1. In contrast, NGC 7469 and IRAS 13120 show variations in the HCN (4-3)/HCN (3-2) ratio with superlinear trends of 1.34 ± 0.05 and 1.3 ± 0.1, respectively. A fit to the pixels from all galaxies in the HCN (4-3)/HCN (3-2) ratio gives a slightly superlinear slope of 1.13 ± 0.03 (Figure 3c).
3.2 Consistency in the CN ratios

The slopes for the CN (2-1)/CN (1-0) ratio are consistent within one standard deviation for all three galaxies in our sample, with an average value of 0.93 ± 0.05 (Figure 3b). We do not resolve any structure in the line ratio maps of individual galaxies (1). There is, however, an offset to lower values for the CN (2-1)/CN (1-0) intensity ratio in NGC 3256 (~ 0.4) compared to NGC 7469 and IRAS 13120 (~ 0.8).

The CN (3-2)/CN (2-1) ratio varies in all the galaxies in our sample (superlinear trends in Table 4 and Figure 3). The slopes of the CN (3-2)/CN (2-1) ratio in NGC 3256 and IRAS 13120 are smaller (1.26 ± 0.05 and 1.11 ± 0.04, respectively) than the strong superlinear trend seen in NGC 7469 (1.55 ± 0.04). We find a higher CN (3-2)/CN (2-1) ratio of 0.45 ± 0.06 in the nucleus of NGC 7469 compared to the non-nuclear value of 0.30 ± 0.04. The nuclear CN (3-2)/CN (2-1) ratio in NGC 7469 more closely resembles the ratio in IRAS 13120 (~ 0.5) than NGC 3256 (~ 0.2). Fitting the pixels from all galaxies, the CN (3-2)/CN (2-1) ratio has a superlinear slope of 1.33 ± 0.02 (Figure 3d). The superlinear trend in this CN (3-2)/CN (2-1) ratio is stronger than the trend in the HCN (4-3)/HCN (3-2) ratio, with values of 1.33 ± 0.02 and 1.13 ± 0.03, respectively.

3.3 Subtle variations in CN/HCN intensity ratios

We find a higher CN (1-0)/HCN (1-0) intensity ratio of 1.7 ± 0.1 in the northern nucleus of NGC 3256, compared to the non-nuclear value of 1.14 ± 0.09. The global CN (1-0)/HCN (1-0) ratio in IRAS 13120 is also 1.7 ± 0.1. The CN (1-0)/HCN (1-0) ratios in IRAS 13120 and in the northern nucleus of NGC 3256 are significantly higher than the ratio of ~ 1.1 in NGC 7469. The southern nucleus and non-nuclear regions of NGC 3256 also have smaller ratios of ~ 1.1.

The variations in the CN (1-0)/HCN (1-0) ratio in NGC 3256 result in a superlinear slope of 1.49 ± 0.06 (Figure 4a). Despite the lack of obvious variations in the different regions of NGC 7469 and IRAS 13120, the slopes indicate subtle variations in the CN (1-0)/HCN (1-0) ratio with values of 1.26 ± 0.05 for NGC 7469 and 0.83 ± 0.05 for IRAS 13120.
For the CN (3-2)/HCN (3-2) ratio, we measure the same ratio of 0.8 ± 0.1 for all regions in NGC 7469 and IRAS 13120. The CN (3-2)/HCN (3-2) ratio in NGC 3256 has lower values between 0.35 – 0.55 (Figure 4b). The lower CN (3-2)/HCN (3-2) ratio in NGC 3256 occurs primarily toward the southern half of the galaxy (Figure 2c). In NGC 3256, the highest CN (3-2)/HCN (3-2) ratio of 0.55 ± 0.08 occurs in the northern nucleus.

4 POTENTIAL PHYSICAL DRIVERS FOR THE OBSERVED RATIO VARIATIONS

4.1 Driving the excitation of HCN

Molecular excitation is dependent on the temperature, density, and optical depth of the molecular gas, as well as the presence of an external radiation source, such as a PDR and/or an XDR (Boger & Sternberg 2005; Meijerink & Spaans 2005; Meijerink et al. 2007). Additionally, molecular excitation can be affected by shocks, which will increase the temperature in the shocked regions of the gas (Martín et al. 2015). The J = 1-0, 3-2 and 4-3 rotational transitions of HCN have energies of 4.25 K, 25.52 K and 42.53 K above the ground state, respectively. HCN excitation to the J = 3-2 and 4-3 transitions thus requires much warmer, dense gas conditions than the J = 1-0 transition. An increase in HCN excitation and intensity ratios is expected in regions surrounding AGN sources (Boger & Sternberg 2005; Meijerink & Spaans 2005; Meijerink et al. 2007; Izumi et al. 2013, 2016). In addition, Saito et al. (2018) find a positive correlation between HCN excitation and SFR surface density when excluding AGN contributions from their analysis. Therefore, we expect to find increased HCN intensity ratios in regions with starburst and/or AGN

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activity. All galaxies in our sample host some combination of AGN and starburst activity that we often cannot separate at our resolutions.

4.1.1 Increased nuclear HCN (3-2)/HCN (1-0) ratio in NGC 7469

The nuclear region of NGC 7469 has an HCN (3-2)/HCN (1-0) ratio of ~ 0.5. There is a circumnuclear disk of cold molecular gas surrounding the AGN of this galaxy and forming stars in the inner ~ 100 pc region (Davies et al. 2004; Izumi et al. 2015). Our observations of NGC 7469 have a resolution of ~ 300 pc and thus blend the influence of both AGN and starburst activity in the circumnuclear disk. This blending makes it challenging to disentangle the effect of the AGN and starburst on the molecular gas, as both will help to increase the HCN (3-2)/HCN (1-0) ratio.

There is evidence of an XDR that is < 50 pc in radius around the AGN in NGC 7469 (Izumi et al. 2015, 2020). This XDR will influence the higher HCN excitation we find in the nuclear region. Izumi et al. (2013, 2016) conclude that there will be high intensities in the HCN sub-mm lines due to IR-pumping in the presence of XDRs. This IR-pumping is a result of X-ray radiation heating dust grains in the vicinity of the XDR and producing strong IR radiation that affects the HCN chemistry. Our future non-LTE analysis with accurate modelling of gas temperatures and densities will help interpret the HCN (4-3)/HCN (3-2) and HCN (3-2)/HCN (1-0) ratios we find in the nucleus of NGC 7469.

Another plausible scenario for the higher HCN line ratios we find around the AGN in NGC 7469 could be shocks from the AGN outflow increasing the temperature of the gas. Shocked gas will have increased temperatures that favour the warm gas chemistry required to produce HCN (Harada et al. 2013; Martin et al. 2015). Martin et al. (2015) find enhanced HCN abundances around the Seyfert 1 AGN in NGC 1097 due to the shocked material produced by the outflows. The peak HCN intensity in their observations is found ~ 200 pc from the galaxy centre, where the X-ray radiation is weaker and the shocks at the base of the outflow are prominent. At our resolution, we cannot distinguish between the effects of the central XDR in NGC 7469 and any shocks in the outflowing gas that could also increase the gas temperature and affect the HCN excitation in this galaxy.

4.1.2 Higher HCN (3-2)/HCN (1-0) ratios in IRAS 13120 than NGC 3256

The HCN (3-2)/HCN (1-0) ratio should increase with increasing starburst contribution due to the positive correlation between HCN excitation and SFR surface density (Saito et al. 2018). Wilson et al. (2019) find ranges of $\Sigma_{\text{SFR}}$ between 10-100 $M_\odot$ yr$^{-1}$ kpc$^{-1}$ for IRAS 13120 compared to 1-25 $M_\odot$ yr$^{-1}$ kpc$^{-1}$ for NGC 3256. Teng et al. (2015) find the 2-10 keV absorption-corrected luminosity of IRAS 13120 is $1.25 \times 10^{45}$ erg s$^{-1}$ and thermal and non-thermal components of the 0.5-2 keV luminosity are consistent with a star formation rate of ~ 170 $M_\odot$ yr$^{-1}$. The lower $\Sigma_{\text{SFR}}$ in NGC 3256 could lead to less HCN excitation (Saito et al. 2018). IRAS 13120 has the highest HCN (3-2)/HCN (1-0) ratio (~ 0.8), while the northern nucleus of NGC 3256 is lower (~ 0.4). We note that the different resolutions in our observations of NGC 3256 (~ 700 pc) and IRAS 13120 (~ 700 pc) mean that we are comparing different spatial scales in the starburst regions of these two galaxies, which might lead to a difference in the HCN (3-2)/HCN (1-0) ratios.

There is no difference in the HCN (3-2)/HCN (1-0) ratio for IRAS 13120 between the global ratio and the continuum peak. Our resolution of ~ 700 pc is insufficient to isolate the nuclear region and so the line ratio will have contributions from both the starburst and AGN components in this galaxy. Veilleux et al. (2013) find that IRAS 13120 has a fast, wide-angle outflow seen in the FIR OH line and suggest an AGN contribution to $L_\text{IR}$ up to 33.4%. This outflow could impact the excitation of HCN through the presence of shocks that could increase the temperature of the gas (Martín et al. 2015).

The southern nucleus of NGC 3256 contains an embedded AGN. Previous work has described this AGN as being inactive or dormant (Sakamoto et al. 2014). The strength of the X-ray emission from this AGN would likely be lower than the other AGN in our sample and we would not expect as significant an increase in the HCN excitation due to X-rays here. Alonso-Herrero et al. (2012) place a 5% upper limit on the contribution from any AGN to the total bolometric luminosity of NGC 3256. Both nuclei in NGC 3256 have been detected in X-rays with Chandra by Lira et al. (2002), although, the authors found no evidence for an AGN in either nucleus. Ohnaka et al. (2015) used Spitzer data and SED fitting to starburst and AGN...
templates to suggest there is an AGN in the southern nucleus with \( L_\nu > 10^{42} \) \( \text{erg s}^{-1} \). The northern nucleus is more likely to host an extreme starburst than an AGN (Neff et al. 2003). These limitations on the AGN in NGC 3256 indicate that the HCN excitation conditions will be almost completely driven by starburst effects on the molecular gas. Any contribution from the AGN in this galaxy to HCN excitation would more likely come from shocks in the outflows from the AGN (e.g. Aalto et al. 2012; Martín et al. 2015).

With higher resolution images and the ability to resolve individual XDRs, we would be able to compare individual ratios based on the relative strengths of the X-rays and outflows created by the AGN.

4.2 Driving the excitation of CN

As with HCN, the excitation conditions of CN will depend on the temperature, density, optical depth, and presence of external radiation fields in the molecular gas. The CN \( N = 1-0, 2-1, \) and 3-2 lines have energies of 5.4 K, 16.3 K, and 32.6 K above ground, respectively. As with HCN, gas traced by higher CN line excitation will be both warmer and/or denser than the \( N = 1-0 \) line.

4.2.1 Superlinear trend in CN (3-2)/CN (2-1)

The superlinear trend we see in the CN (3-2)/CN (2-1) intensity ratio indicates that the different global ratios between our galaxies are correlated with increased intensity of line emission. We observe higher intensities on average in IRAS 13120 than in NGC 3256 and NGC 7469 for both CN lines. These higher ratios could indicate molecular gas that is more dense with higher temperatures, leading to higher CN excitation. Additionally, because IRAS 13120 is a ULIRG with a high SFR surface density, we could expect a higher external UV field interacting with the molecular gas and leading to an increase in CN abundance by activating the CN formation pathways (i.e. HCN photodissociation and ionization of carbon; Boger & Sternberg 2005).

In NGC 7469, the CN (3-2)/CN (2-1) ratio appears to be lowest toward the edges of the map, with mid-range values in the gas surrounding the nucleus, and the highest ratio in the nuclear pixels. We find that the CN (3-2)/CN (2-1) ratio increases from the non-nuclear value of 0.30 ± 0.04 to a value of 0.45 ± 0.06 in the nucleus. The increased CN excitation in this galaxy can thus be attributed to the conditions near the central AGN and surrounding starburst ring. It is apparent that the environment in the nucleus of NGC 7469 has a significant impact on the higher CN excitation. X-rays from the AGN will enhance the bulk temperature of the gas, while PDRs will have higher surface temperatures. In this way, the heating of the gas in the centre of NGC 7469 could be reflected in the CN excitation. We note that it is unclear how far the X-rays will reach, however, and that at our resolution of ~ 300 pc in this galaxy we do not resolve the XDR to localize the extent of AGN gas heating.

While we do not find any localized enhancement in the CN (3-2)/CN (2-1) ratio in IRAS 13120, the CN (3-2) emission is high throughout the galaxy. We find higher CN (3-2)/CN (2-1) ratios (> 0.5) in both the global and continuum peak measurements of IRAS 13120 than in any regions of NGC 3256 (< 0.3) or the non-nuclear region of NGC 7469 (~ 0.3). The nuclear region of NGC 7469 has a CN (3-2)/CN (2-1) ratio of 0.45, which is similar to the global ratio in IRAS 13120. Our ability to resolve different contributions to the excitation to CN (3-2) in IRAS 13120 is limited, while we are better able to localize the effect of the AGN on this ratio in NGC 7469. In any case, both the starburst and AGN contributions in IRAS 13120 will help to increase the excitation of CN to higher values than we find in NGC 3256 and the non-nuclear region of NGC 7469.

We also find a local increase in the CN (3-2)/CN (2-1) ratio in the northern nucleus of NGC 3256 (~ 0.27) compared to the southern nucleus (~ 0.22) and non-nuclear pixels (~ 0.21). However, these ratios are not significantly different once uncertainties on the measured values are taken into account. The higher CN (3-2)/CN (2-1) ratio in the northern nucleus of NGC 3256 could be caused by the more widespread starburst in the northern nucleus. At our resolution (~ 470 pc), we do not completely resolve the edge-on southern nucleus and therefore will have lower excitation gas from non-nuclear regions mixed into the beam. Sakamoto et al. (2014) also observed enhanced CN emission associated with the molecular outflows in the northern nucleus of NGC 3256. The CN emission could thus be enhanced as a result of outflow chemistry that is influenced by far-UV emission from massive stars that are formed in the outflow (e.g. Cicone et al. 2020). Cicone et al. (2020) note, however, that CN enhancement in the outflow of NGC 3256 is weaker than in Mrk 231.

4.2.2 The CN (2-1)/CN (1-0) intensity ratio

The CN (2-1)/CN (1-0) ratio has a consistent slope (~ 0.9) and ratio (~ 0.8) in NGC 7469 and IRAS 13120. This result suggests that the conditions in these two galaxies produce similar excitation conditions for CN. Radiation in the nuclear region of NGC 7469 is dominated by X-rays from the AGN. However, at our resolution of ~ 304 pc, we will include more than just the X-ray effects on the circumnuclear disk of molecular gas contained within the central beam. The 2-10 keV luminosity of the type 1 Seyfert nucleus is \( L_{2-10 \text{ keV}} = 1.5 \times 10^{41} \text{ erg s}^{-1} \) (Liu et al. 2014). The AGN accounts for 30 – 40% of the IR luminosity in NGC 7469 (Alonso-Herrero et al. 2012). IRAS 13120 has a Compton thick AGN that accounts for ~ 33.4% of the total IR luminosity (Iwasawa et al. 2011; Teng et al. 2015; Veilleux et al. 2013). Teng et al. (2015) find the 2-10 keV absorption-corrected luminosity of IRAS 13120 is 1.25 × 10^{42} \text{ erg s}^{-1}, which is comparable to that in NGC 7469. At our resolution, we are unable to distinguish between effects from the AGN and/or starburst activity in IRAS 13120. The presence of both UV and X-ray radiation fields should increase the CN abundance (Boger & Sternberg 2005; Meijerink & Spaans 2005; Meijerink et al. 2007).

Our results indicate that the different radiation fields in NGC 7469 and IRAS 13120 seem to affect the molecular gas in a similar way.

The offset to lower CN (2-1)/CN (1-0) ratios on the whole in NGC 3256 is interesting. The CN (1-0) data for NGC 3256 is our only data set from ALMA Cycle 0. As such, the flux measurements and intensity ratios using this line in this galaxy have the highest calibration uncertainties of our sample. Given the similarity of the slopes in the CN (2-1)/CN (1-0) ratio of all three galaxies, the offset in the ratio in NGC 3256 could be due to large observational uncertainties. The difference in the CN (2-1)/CN (1-0) ratio in NGC 3256 of ~ 0.45 compared to the ratios in NGC 7469 and IRAS 13120 of ~ 0.8 (Figure 3) should be treated with caution until this is confirmed with follow-up observations. If real, the offset to lower ratios could be due to different excitation conditions in the dense gas of this galaxy. The small contribution from any AGN in NGC 3256 (< 5%: Alonso-Herrero et al. 2012) could perhaps contribute to the lower (~ 0.5) CN (2-1)/CN (1-0) ratio compared to those in NGC 7469 and IRAS 13120 (both > 0.75). The X-ray luminosity in NGC 3256 is \( L_{2-10 \text{ keV}} = 1.5 \times 10^{40} \text{ erg s}^{-1} \) (Ohyama et al. 2015), three orders of magnitude smaller than both NGC 7469 and IRAS 13120. The models by Meijerink et al. (2007) indicate that CN abundance will be more strongly affected in the presence of XDRs than PDRs.
suggesting that the higher CN intensity ratios we find in IRAS 13120 and NGC 7469 could be a result of the influence of the AGN. If both the UV and X-ray fields created in NGC 3256 are weaker than those found in IRAS 13120 or NGC 7469, the CN (1-0) emission could dominate over the CN (2-1) as conditions are not favourable for CN to be excited to higher levels. Future non-LTE modelling to estimate the opacities, temperatures, and densities in these galaxies will allow us to better explain the offset we see in the CN (2-1)/CN (1-0) ratio in NGC 3256.

4.3 The CN/HCN intensity ratio as a probe of radiation field

The HCN J = 1-0 and CN N = 1-0 transitions have energies of ~ 4.25 K and ~ 5.4 K, respectively, suggesting they trace molecular gas of similar temperatures and densities. In fact, recent work using ALMA observations has found that the global observed CN (1-0)/HCN (1-0) intensity ratio is nearly constant over 3 orders of magnitude change in similar temperatures and densities. In fact, recent work using ALMA (Wilson et al. in prep). The average CN (1-0)/HCN (1-0) ratio in their sample of 9 galaxies is 0.866 ± 0.096, with a standard deviation of 0.29. Wilson et al. (in prep) use an additional criteria of detecting 93 GHz radio continuum emission in order to compare their CN (1-0)/HCN (1-0) ratios to $\Sigma_{\text{SFR}}$. Since we are not requiring 93 GHz continuum detections, we can extend our analysis to the lower line intensities where we find most of our variations in the CN (1-0)/HCN (1-0) ratios. Including both CN N = 1-0 hyperfine lines, we find ratios of 1.1 – 1.7 for the CN (1-0)/HCN (1-0) ratio.

We find variations in the observed intensity ratios between CN and HCN lines, particularly in regions associated with high starburst activity. For example, the CN (1-0)/HCN (1-0) ratio in IRAS 13120 has a value of 1.7 ± 0.1. IRAS 13120 is a ULIRG with a higher $\Sigma_{\text{SFR}}$ than both NGC 3256 and NGC 7469. The increased star formation activity in this galaxy should be indicative of an increase in UV radiation due to young, massive star formation, which could indicate that the UV radiation more strongly affects the gas properties. The northern nucleus of NGC 3256 has the same CN (1-0)/HCN (1-0) ratio of 1.7±0.1 and is also a starburst with a high UV field (Sakamoto et al. 2014). These results will be explored further using non-LTE analysis to calculate the CN/HCN abundance ratio in the presence of UV fields using a PDR model in a future paper.

The CN (1-0)/HCN (1-0) ratio in the southern nucleus of NGC 3256 (~ 1.1) is more comparable with that of NGC 7469 (~ 1.1). In the southern nucleus of NGC 3256, the AGN and disk are more compact and therefore the sphere of influence of any enhancement from a UV field due to star formation or X-rays from the AGN are relatively small compared to the extent of the starburst in the northern nucleus. The exact impact of X-ray versus UV radiation fields on CN and HCN excitation is beyond the scope of this work.

We find a relatively constant CN (1-0)/HCN (1-0) intensity ratio across NGC 7469, suggesting that the presence of an AGN and XDR does not affect the ratio as significantly as the starbursts do in IRAS 13120 and the northern nucleus of NGC 3256. Non-LTE modelling to determine the CN and HCN abundance ratio is needed to determine if the CN/HCN abundance ratio is also constant. Izumi et al. (2015) suggest that high-temperature gas-phase chemistry in NGC 7469 could significantly enhance the abundance of HCN through an activated formation pathway from CN when T > 300 K (e.g. Harada et al. 2010).

Throughout this paper, we have described the CN/HCN intensity ratios from observations of multiple CN and HCN lines in three galaxies. We find variations in the CN (1-0)/HCN (1-0) intensity ratio in NGC 3256, from values of ~ 1 in the non-nuclear pixels to values of ~ 1.7 in the northern nucleus. A simple LTE calculation suggests that the CN and HCN column densities should be in the range of $10^{13} - 10^{17}$ cm$^{-2}$. We used this column density range to perform a simple non-LTE calculation with the online version of RADEX$^7$ (van der Tak et al. 2007). Using a gas temperature of 40 K, an H$_2$ number density of $10^7$ cm$^{-3}$, and a line width of 10 km s$^{-1}$, we reproduce CN (1-0)/HCN (1-0) ~ 1 in the non-nuclear pixels of NGC 3256 with a column density ratio of $N_{\text{CN}}/N_{\text{HCN}}$ ~ 10. The intensity ratio of CN (1-0)/HCN (1-0) ~ 1.7 in the northern nucleus of NGC 326 is reproduced with $N_{\text{CN}}/N_{\text{HCN}}$ ~ 24. Further non-LTE analysis will be performed in future work, with more analysis on the specific gas properties and optical depths in NGC 3256, NGC 7469, and IRAS 13120 that can produce our observed CN and HCN intensity ratios. In this way, we will be able to compare our data to the expectation that the CN/HCN abundance ratio should increase in the presence of a starburst or increased radiation field (c.f. Fuente et al. 1993, 1995; Aalto et al. 2002; Boger & Sternberg 2005; Meijerink & Spaans 2005).

5 CONCLUSIONS

In our observational study of CN and HCN lines in a sample of three U/LIRG galaxies, we find subtle variations in line intensity ratios within and among galaxies. These variations can be localized to regions of starburst and/or AGN activity. Both CN and HCN excitation are enhanced in the higher intensity regions.

Our main conclusions are summarized as follows:

(i) We find variations in the HCN (3-2)/HCN (1-0) ratio in NGC 3256 and NGC 7469. Both galaxies have superlinear trends with slopes > 1.5. The global and non-nuclear HCN (3-2)/HCN (1-0) and HCN (4-3)/HCN (1-0) ratios are the same (~ 0.3) in the two galaxies, which suggests similar excitation conditions. The enhanced HCN excitation is seen in both nuclei of NGC 3256 and in the nucleus of NGC 7469. The HCN (3-2)/HCN (1-0) ratio in NGC 7469 shows the largest difference between non-nuclear (0.33 ± 0.04) and nuclear (0.52 ± 0.06) regions. We attribute the regions of enhanced HCN excitation in NGC 7469 to the central AGN, where both an XDR and shocks in outflows can heat the gas and increase HCN excitation. The HCN (3-2)/HCN (1-0) ratios are increased relative to the disk in both nuclei of NGC 3256 due to both starburst and AGN effects, although to a lesser extent than in the nucleus of NGC 7469.

(ii) The HCN (3-2)/HCN (1-0) ratio in IRAS 13120 is higher (~ 0.8) than both NGC 3256 (~ 0.27 – 0.39) and NGC 7469 (~ 0.33 – 0.52). We find little spatial variation in the HCN excitation of IRAS 13120. We attribute the higher HCN (3-2)/HCN (1-0) ratios to the higher star formation rate surface density and contribution of the Compton-thick AGN to the total IR luminosity.

(iii) We find relatively constant CN (2-1)/CN (1-0) ratios and subtle variations in the CN (3-2)/CN (2-1) ratio in our galaxy sample. We find an offset to lower ratios of ~ 0.45 in the CN (2-1)/CN (1-0) ratio in NGC 3256 compared to ~ 0.8 in NGC 7469 and IRAS 13120.

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6 In their analysis, which includes the galaxies NGC 3256, NGC 7469, and IRAS 13120, they consider only the strong CN (1-0) hyperfine line, which accounts for ~ 70% of the total CN (1-0) line emission for optically thin lines (Shirley 2015; Meier et al. 2015).

7 https://home.strw.leidenuniv.nl/moldata/radex.html
We note that this offset could arise from the difference in the fraction of AGN contribution to the total IR luminosity in NGC 3256 (< 5%) compared to NGC 7469 (32 – 40%) and IRAS 13120 (17 – 33.4%). Further non-LTE modelling analysis may help explain the origin of this offset.

(iv) The CN (1-0)/HCN (1-0) ratio is higher (1.7) in the northern nucleus of NGC 3256 than in the rest of the galaxy. We find a similarly high CN (1-0)/HCN (1-0) ratio in the starburst galaxy IRAS 13120. The intense starburst activity in the regions with large CN (1-0)/HCN (1-0) ratios suggests the enhanced ratio may be due to increased exposure to UV radiation fields. We note that opacity effects could potentially be important and determining the CN/HCN abundance ratios is the next step in this analysis. The CN (1-0)/HCN (1-0) ratio in NGC 7469 is ~ 1.1 and is relatively constant across the entire galaxy. NGC 7469 is dominated by a nuclear AGN, which does not appear to affect the CN/HCN ratio as significantly as the starbursts in the other galaxies.

The more detailed examination of CN and HCN conducted in this work helps confirm the conclusion of Wilson et al. (in prep.) that CN is a viable tracer of dense gas in galaxies. However, it is clear that more work is needed to understand the complex relationship between CN and HCN in dense gas exposed to strong UV and X-ray radiation fields. Future non-LTE radiative transfer modelling of the CN and HCN lines in this galaxy sample will allow us to determine physical properties such as gas temperature, density, opacity, and molecular abundance. Using the molecular abundances, we will be able to compare our results to models predicting CN/HCN abundance ratios and further explore the effects of radiation fields on molecular gas properties.

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ADS/JAO.ALMA#2012.1.00165.S
ADS/JAO.ALMA#2013.1.00218.S
ADS/JAO.ALMA#2013.1.00379.S
ADS/JAO.ALMA#2015.1.00102.S
ADS/JAO.ALMA#2015.1.00287.S
ADS/JAO.ALMA#2015.1.00412.S
ADS/JAO.ALMA#2015.1.00993.S
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DATA AVAILABILITY

This research is based on public data available in the ALMA archive and the exact project IDs can be found in Table 2. The whole dataset and processed data cubes can be obtained from the authors by request.

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\[8 \text{ https://linmix.readthedocs.io/en/latest/src/linmix.html}\]
A2 Full measured intensity ratios

Table A3 presents the measured CN and HCN intensity ratios in the global, nuclear, and non-nuclear regions of NGC 3256, NGC 7469, and IRAS 13210 for additional pairs of lines beyond those given in Table 4.

A3 Additional pixel-by-pixel comparisons

We present additional pixel-by-pixel correlations between HCN (4-3) and HCN (1-0), CN (3-2) and CN (1-0), CN (2-1) and HCN (3-2), and CN (3-2) and HCN (4-3) using the intensities in the re-sampled moment 0 maps in the scatter plots of Figures A5 and A4. In Table A2, we present the measured slopes and intercepts from the Linmix linear regression fits to the various line ratios explored in the pixel-by-pixel comparisons. In this table we also include the HCN (4-3)/HCN (1-0), CN (3-2)/CN (1-0), CN (2-1)/HCN (3-2), and CN (3-2)/HCN (4-3) ratios. The linear fits were in the form $\log(y) = m \log(x) + b$ on to the pixel-by-pixel comparisons of the re-sampled moment 0 maps and we present both $m$ and $b$ values here. Inclination corrections are as described in Section 2.4. Fits to all pixels from all galaxies for the various line ratios are presented in Table A2, as well.

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Table A1. Data reduction imaging properties and line sensitivities.

| Imaging property | NGC 3256 | NGC 7469 | IRAS 13120 |
|------------------|----------|----------|------------|
| Beam size \(a\) | 2.2''    | 0.95''   | 1.1''      |
| Velocity resolution (\(\text{km s}^{-1}\) \(b\)) | 26.43    | 20.68    | 20.64      |
| Moment 0 map velocity ranges (\(\text{km s}^{-1}\)) |          |          |            |
| HCN(1-0)        | 2584.30 – 2980.75 | 4632.72 – 5087.68 | 8744.48 – 9363.68 |
| HCN(3-2)        | 2584.30 – 2980.75 | 4624.08 – 5079.04 | 8744.48 – 9363.68 |
| HCN(4-3)        | 2584.30 – 2980.75 | 4632.72 – 5087.68 | 8744.48 – 9363.68 |
| CN(1-0)         | 2555.72 – 3850.79 | 4503.40 – 5087.68 | 8744.48 – 9363.68 |
| CN(2-1)         | 1805.72 – 3021.50 | 3865.44 – 4940.80 | 8003.20 – 9262.24 |
| CN(3-2)         | 1905.72 – 3121.50 | 3944.76 – 5082.16 | 8044.48 – 9179.68 |
| Smoothed cube line sensitivities (mJy beam\(^{-1}\)) | \(c\) | | |
| HCN(1-0)        | 0.241     | 0.280    | 0.958      |
| HCN(3-2)        | 1.624     | 0.478    | 0.628      |
| HCN(4-3)        | 1.059     | 0.872    | 2.166      |
| CN(1-0)         | 0.635     | 0.508    | 1.220      |
| CN(2-1)         | 1.021     | 0.277    | 0.424      |
| CN(3-2)         | 0.570     | 0.753    | 1.035      |

Notes: 
\(a\) The beams were smoothed, rounded, and matched to this resolution for all lines.
\(b\) The velocity resolution and pixel size were matched between all lines.
\(c\) The smoothed cube line sensitivities were determined using the same line-free channels as for the dirty cubes. Integrated intensity map 3.5\(\sigma\) cut-offs used these smoothed cube sensitivities.

Figure A1. Integrated intensity maps in K \(\text{km s}^{-1}\) of NGC 3256 for all six lines. In the lower left corner, the circle indicates the beam size of 2.2'' and the scale bar is set at 500 pc. NGC 3256 clearly has a peak intensity in the northern nucleus. The southern nucleus also shows a local increase in intensity. One spiral arm feature can be seen at our sensitivity in the top left corner of the maps.
Figure A2. Integrated intensity maps in K km s$^{-1}$ of NGC 7469 for all six lines. In the lower left corner, the circle indicates the beam size of 0.95″ and the scale bar is set at 500 pc. The peak intensity corresponds to the central nuclear region hosting the AGN. We see some increase in intensity in the region surrounding the nucleus, potentially corresponding to the starburst ring found in this galaxy.

Figure A3. Integrated intensity maps in K km s$^{-1}$ of IRAS 13120 for all six lines. In the lower left corner, the circle indicates the beam size of 1.1″ and the scale bar is set at 500 pc. The furthest galaxy in our sample, IRAS 13120 has quite a symmetric appearance. The intensities of individual lines at the central peak are much higher than those detected in NGC 3256 and NGC 7469.
**Figure A4.** Pixel-by-pixel comparisons of individual line intensities in each galaxy. The higher-J/N transition of the two lines is always plotted on the y-axis. The black dotted line is the one-to-one line. Blue squares correspond to NGC 3256; pink inverted triangles correspond to NGC 7469; purple circles correspond to IRAS 13120. The dashed lines in each colour represent the Linmix fits for each galaxy, with the shaded region representing the 95% confidence interval of these fits. Large stars show the global ratio in each galaxy and large pentagons show the non-nuclear ratios in NGC 3256 and NGC 7469. Large triangles show ratios for the northern nucleus of NGC 3256, the nucleus of NGC 7649, and the continuum peak in IRAS 13120, while the large inverted triangle shows the ratio for the southern nucleus of NGC 3256. The black dot-dashed lines and shaded regions in (c) and (d) show Linmix fits to all three galaxies.

**Figure A5.** Pixel-by-pixel comparisons for CN compared to HCN for selected transitions. For plot descriptions, refer to the caption of Figure A4.
Table A2. Slope and intercept from pixel-by-pixel comparisons.

| Ratio     | Params. | NGC 3256 | NGC 7469 | IRAS 13120 | All pixels |
|-----------|---------|----------|----------|------------|------------|
| HCN (3-2) |         |          |          |            |            |
| HCN (1-0) | m       | 1.53 ± 0.07 | 1.55 ± 0.05 | 0.84 ± 0.06 | 1.46 ± 0.04 |
|           | b       | −1.4 ± 0.1  | −1.45 ± 0.08 | 0.3 ± 0.1   | −1.26 ± 0.07 |
| HCN (4-3) |         |          |          |            |            |
| HCN (1-0) | m       | 1.50 ± 0.08 | 2.1 ± 0.1  | 1.19 ± 0.07 | 1.46 ± 0.04 |
|           | b       | −1.7 ± 0.1  | −2.9 ± 0.2  | −0.7 ± 0.1  | −1.26 ± 0.07 |
| HCN (4-3) |         |          |          |            |            |
| HCN (3-2) | m       | 0.99 ± 0.04 | 1.34 ± 0.05 | 1.3 ± 0.1   | 1.13 ± 0.03 |
|           | b       | −0.34 ± 0.04 | −0.91 ± 0.06 | −0.9 ± 0.2  | −0.55 ± 0.04 |
| CN (2-1)  |         |          |          |            |            |
| CN (1-0)  | m       | 0.95 ± 0.03 | 0.93 ± 0.03 | 0.91 ± 0.02 | 1.07 ± 0.02 |
|           | b       | −0.29 ± 0.05 | −0.00 ± 0.06 | 0.11 ± 0.06 | −0.33 ± 0.04 |
| CN (3-2)  |         |          |          |            |            |
| CN (1-0)  | m       | 1.16 ± 0.07 | 1.42 ± 0.07 | 0.99 ± 0.05 | 1.07 ± 0.02 |
|           | b       | −1.4 ± 0.1  | −1.4 ± 0.1  | −0.4 ± 0.1  | −0.33 ± 0.04 |
| CN (3-2)  |         |          |          |            |            |
| CN (2-1)  | m       | 1.26 ± 0.05 | 1.55 ± 0.04 | 1.11 ± 0.04 | 1.33 ± 0.02 |
|           | b       | −1.05 ± 0.06 | −1.47 ± 0.06 | −0.56 ± 0.08 | −1.11 ± 0.03 |
| CN (1-0)  |         |          |          |            |            |
| HCN (1-0) | m       | 1.49 ± 0.06 | 1.26 ± 0.05 | 0.83 ± 0.05 | 1.25 ± 0.03 |
|           | b       | −0.74 ± 0.09 | −0.40 ± 0.08 | 0.6 ± 0.1  | −0.38 ± 0.06 |
| CN (2-1)  |         |          |          |            |            |
| HCN (3-2) | m       | 0.89 ± 0.05 | 0.77 ± 0.02 | 0.86 ± 0.04 | 0.91 ± 0.02 |
|           | b       | 0.35 ± 0.05 | 0.70 ± 0.03 | 0.54 ± 0.08 | 0.44 ± 0.03 |
| CN (3-2)  |         |          |          |            |            |
| HCN (2-1) | m       | 1.06 ± 0.09 | 1.17 ± 0.05 | 0.94 ± 0.06 | 1.21 ± 0.04 |
|           | b       | −0.56 ± 0.08 | −0.35 ± 0.06 | 0.1 ± 0.1  | −0.52 ± 0.05 |
| CN (3-2)  |         |          |          |            |            |
| HCN (3-3) | m       | 1.14 ± 0.07 | 0.85 ± 0.03 | 0.69 ± 0.04 | 1.01 ± 0.04 |
|           | b       | −0.23 ± 0.05 | 0.46 ± 0.03 | 0.75 ± 0.07 | 0.12 ± 0.04 |

Notes: aFit parameters are determined from Linmix linear fitting of the form log(y) = m log(x) + b.
### Table A3. Additional CN and HCN measured intensity ratios.

| Region          | HCN (3-2) | HCN (4-3) | HCN (4-3) | CN (2-1) | CN (3-2) | CN (3-2) | CN (3-2) | CN (3-2) |
|-----------------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|
|                 | HCN (1-0) | HCN (1-0) | HCN (1-0) | CN (1-0) | CN (2-1) | CN (2-1) | CN (2-1) | CN (2-1) |
| NGC 3256        |           |           |           |           |           |           |           |           |
| Global          | 0.29(3)   | 0.13(2)   | 0.46(7)   | 0.44(5)   | 0.10(1)   | 0.22(3)   | 1.24(9)   | 1.9(3)    |
| North nucleus   | 0.39(4)   | 0.18(2)   | 0.48(7)   | 0.47(5)   | 0.13(1)   | 0.27(4)   | 1.7(1)    | 2.1(3)    |
| South nucleus   | 0.33(4)   | 0.15(2)   | 0.46(7)   | 0.44(5)   | 0.10(1)   | 0.22(3)   | 1.12(9)   | 1.5(2)    |
| Non-nuclear     | 0.27(3)   | 0.12(1)   | 0.45(6)   | 0.43(5)   | 0.09(1)   | 0.21(3)   | 1.14(8)   | 1.8(3)    |
| NGC 7469        |           |           |           |           |           |           |           |           |
| Global          | 0.36(4)   | 0.14(2)   | 0.40(6)   | 0.76(9)   | 0.25(3)   | 0.32(5)   | 1.12(8)   | 2.4(3)    |
| Nucleus         | 0.52(6)   | 0.25(3)   | 0.48(7)   | 0.80(9)   | 0.36(4)   | 0.45(6)   | 1.13(9)   | 1.7(2)    |
| Non-nuclear     | 0.33(4)   | 0.12(1)   | 0.37(5)   | 0.76(9)   | 0.22(3)   | 0.30(4)   | 1.11(8)   | 2.6(4)    |
| IRAS 13120      |           |           |           |           |           |           |           |           |
| Global          | 0.80(9)   | 0.61(7)   | 0.8(1)    | 0.77(9)   | 0.40(5)   | 0.53(8)   | 1.7(1)    | 1.6(2)    |
| Continuum peak  | 0.79(9)   | 0.70(8)   | 0.9(1)    | 0.80(9)   | 0.45(5)   | 0.57(8)   | 1.5(1)    | 1.5(2)    |

**Notes:** The uncertainties presented here are the measurement and calibration uncertainties and given as the uncertainty on the last digit (i.e. $0.29(3) = 0.29 \pm 0.03$). Ratios include the 5% (Band 3) and 10% (Band 6 and 7) ALMA calibration uncertainties for each line.