FROM PRIMITIVE FORM TO MIRROR SYMMETRY

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ABSTRACT. This is a report on the recent joint work [29] on LG-LG mirror symmetry for the 14 exceptional unimodular singularities.¹

1. INTRODUCTION

In the early 80’s, the author introduced the theory of primitive forms [38, 40, 42], which studies the period integrals of a primitive form over cycles which are vanishing to isolated critical points of a function. One important consequence of the theory is that a primitive form induces a flat structure [42] on the deformation parameter space of the function, including the flat metric, the flat coordinate system and the potential function (which is later called the prepotential). Later on, in the early 90’s, Dubrovin [12] studied the 2D topological field theory of genus zero curves and found the same structure, which is axiomatized to, so-called, the Frobenius manifolds structures. Lots of Frobenius manifolds structures are found. They include the examples constructed on the orbit spaces of Coxeter groups [39, 45], [13], [48] (which contain the first cases found before the primitive form theory), the one constructed by primitive forms [42], [10], [11], [23], [50], the quantum cohomology rings, Barannikov-Kontsevich construction [41] using polyvector fields of a Calabi-Yau manifold, and the FJRW-theory (the A-model of quantum singularity theory) [14, 15].

The Gromov-Witten theory [25] counts pseudoholomorphic curves in a given symplectic manifold. Its application to the symplectic structure on a Kähler manifold was extensively studied from a view point of the mirror symmetry. Here the mirror symmetry is one of the dualities in physics and had a strong impact on mathematics [21, 22]. Namely, it asserts that certain data counted from symplectic geometry (the A-model side) should be equivalent to that from the complex structure of the mirror manifold (the B-model side). There are several different levels of formulation of the mirror symmetry such as the categorical level [16, 24], the geometric level [49], or the equivalence of the genus zero theory on the A-model side with the variation of Hodge structures on the B-model side [19, 30].

¹ Present note is worked out with the help of the coauthors Changzheng Li, Si Li and Yefeng Shen, to whom the author expresses his deep gratitudes.
Primitive forms are about universal deformations $F$ of functions, giving flat structures on the deformation spaces. Hence, the theory is relevant in the complex geometric (B-model) aspects of $N=(2,2)$ supersymmetric Landau-Ginzburg (LG) theory with the superpotential $F$. However, this pattern of the mirror symmetry was not mathematically rigorously worked out until recently. This is because of (1) lack of mathematical theory of A-model LG-theory at that early time, and (2) the difficulty of calculating primitive forms until recently, where explicit expressions of primitive forms were known only for weighted homogeneous polynomials of central charge less than or equal to 1. Both difficulties were resolved as follows.

In around 2007, Fan, Jarvis and Ruan constructed a so-called quantum singularity theory by counting virtual cycles associated with a weighted homogeneous polynomial, whose potential (generating series) gives again a Frobenius manifold structure on the so-called FJRW state space \cite{14,15}. This is considered as an A-model Landau-Ginzburg theory. They immediately realized that such Frobenius manifold for an ADE-polynomial $W$ is actually isomorphic to the Frobenius manifold arising from the primitive form (i.e. B-side) of another ADE-polynomial $W^T$. The superpotential polynomial $W^T$ on B-side is obtained by the transposition of exponents of monomials in the polynomial $W$ on A-side \cite{2,3,26}, which is later called Berglund-Hübsch-Krawitz mirror. As an application of this mirror theorem, they solved the so-called generalized Witten conjecture, which says that the generating functions arising from the Landau-Ginzburg model for ADE-singularities should be governed by some ADE-integrable hierarchies. Also a similar observation for simply elliptic singularities \cite{37} was achieved \cite{27,32,33}. We remark that the relationship between FJRW theory and Gromov-Witten theory (both are in A-model side) is studied under the name of LG/CY-correspondence, for which one is referred to \cite{6,7,27,32,35}.

On the other hand, in 2013, jointly with Changzheng Li and Si Li, the author came to a new perturbative construction of primitive forms \cite{28}, where Birkhoff decomposition theorem used in the original formulation \cite{42} was replaced by the asymptotic expansion of oscillatory integrals. This enables us to calculate primitive forms explicitly as a power series in an algorithmic way (at least for weighted homogeneous polynomials). With the perturbative approach, we can calculate further the flat coordinate system and the pre-potential function up to any finite order. This will be sufficient to determine the flat coordinate system and the pre-potential function with a help of WDVV-equations.

These two new developments thoroughly changed the view on the LG-LG mirror symmetry. Namely, up to a choice of primitive forms, one asks whether the pre-potential attached to FJRW theory for a weighted homogeneous polynomial could coincide with the prepotential associated to a primitive form for the mirror dual-polynomial. Such a mirror
symmetry is called the Landau-Ginzburg to Landau-Ginzburg (LG-LG) mirror symmetry. Its study has developed rapidly in the last years. In the present note, we briefly introduce the theories on both sides and the mirror map construction connecting them. Then, we confirm the LG-LG mirror symmetry for the 14 unimodular singularities, which are the first case of weighted homogeneous polynomials whose central charge exceeds 1.

**Remark.** We remark that the primitive form theory depends only on the analytic equivalence class of the singularity of the function $W^T$, although associated primitive forms may not be unique but form a family. On the other hand, FJRW theory depends on the polynomial $W$ itself together with a symmetry group of $W$. Hence, to achieve the mirror symmetry to FJRW theory on $W$, both the analytic equivalence class of $W^T$ and the choice of the primitive form for $W^T$ depend on the choice of the polynomial $W$. We do not yet understand this phenomenon conceptually (c.f. [29] Remark 4.9. (2)).

2. **Primitve Form Theory.**

The origin of a Landau-Ginzburg B-model (with respect to trivial group symmetry) at genus zero is the theory of primitive forms [28, 38, 40–42, 44]. The starting data of the theory is a holomorphic function $f : (X, 0) \rightarrow (\mathbb{C}, 0)$ defined on a Stein domain $X \subset \mathbb{C}^n$ with finite critical points. For our purpose on the LG-LG mirror symmetry, it is sufficient to consider a weighted homogeneous polynomial $f = f(x_1, \ldots, x_n)$ with an isolated critical point at the origin $0 \in X = \mathbb{C}^n$,

$$f(\lambda^{q_1}x_1, \ldots, \lambda^{q_n}x_n) = \lambda f(x_1, \ldots, x_n), \ \forall \lambda \in \mathbb{C}^*,$$

Here $(q_1, \ldots, q_n)$ in $\mathbb{Q}_{\geq 0}^n$ are called the weights of the coordinates $(x_1, \ldots, x_n)$, and each weight $0 < q_i \leq \frac{1}{2}$ is unique [36]. In [41], the author introduced the formal completion of the Brieskorn lattice together with a semi-infinite $z$-adic filtration by a formal variable $z$:

$$\hat{H}_f(0) : = \Omega^n_{X, 0}[[z]]/(df + zd)\Omega^{n-1}_{X, 0}[[z]],$$

and constructed a higher residue pairing

$$K_f : \hat{H}_f(0) \otimes \hat{H}_f(0) \rightarrow z^n\mathbb{C}[[z]]$$

which satisfies a number of properties, and plays a key role in the theory of primitive forms. A universal unfolding of $f$ is given by

$$F : (X \times S, 0 \times 0) \rightarrow (\mathbb{C}, 0), \quad F(x, s) = f(x) + \sum_{\alpha=1}^{\mu} s_\alpha \phi_\alpha,$$
where \( \{ \phi_{1}, \cdots, \phi_{\mu} \} \subset \mathbb{C}[x] \) are weighted homogeneous polynomials representing an additive basis of the Jacobian algebra \( \text{Jac}(f) \), and \( s = \{ s_{\alpha} \}_{\alpha=1, \cdots, \mu} \) parametrizes the deformation space \( S \subset \mathbb{C}^{\mu} \). Using, so called, Kodaira-Spencer map: \( \sum_{i} a_{i} \partial x_{i} \mapsto \sum_{i} a_{i} \phi_{i} \), the tangent bundle of \( S \) is identified with the Jacobi ring of \( F \), which gives a ring structure (Frobenius algebra structure) and a natural inner product \( \mathcal{I} \) (the first residue pairing) on the tangent bundle of \( S \). There is a family version \( \hat{H}_{F}^{(0)} \) (resp. \( K_{F} \)) of \( \hat{H}_{f}^{(0)} \) (resp. \( K_{f} \)) with respect to the universal unfolding \( F \). We remark that in the recent work \cite{28} by C. Li, S. Li and the author, an alternate complex differential geometric approach to the module \( \hat{H}_{F}^{(0)} \) is developed. Therein we give a simple construction of the higher residue pairing by using integration of compactly supported polyvector fields.

A primitive form is a section \( \zeta \in \Gamma(S, \hat{H}_{F}^{(0)}) \), represented by a relative holomorphic volume form \( \zeta = P(x, s) d^{n}x \) \( (d^{n}x = dx_{1} \cdots dx_{n}) \) on \( X \times S \), satisfying the properties of primitivity, orthogonality, holonomicity, and homogeneity, described by bilinear equations on \( \zeta \) using the higher residue pairing \( K_{F} \) together with Gauss-Manin connection on \( \hat{H}_{F}^{(0)} \).

Roughly speaking, the submodule of \( \hat{H}_{F}^{(0)} \) consisting of the covariant differentiations of a primitive form by the tangent vectors of \( S \) forms a splitting factor to the adic filtration on \( \hat{H}_{F}^{(0)} \) defined by (the powers of) \( z \), i.e. \( \hat{H}_{F}^{(0)} \simeq \mathcal{T}_{S} \oplus z \cdot \hat{H}_{F}^{(0)} \). In this way, properties of the primitive form are transferred to the splitting factor i.e. to the tangent bundle of \( S \), and, hence, the space \( S \) obtains a differential geometric structure, called the flat structure (= the Frobenius manifold structure) associated with \( \zeta \). For instance, the orthogonality property of \( \zeta \) gives a flat metric \( J \) (i.e. the first residue pairing) on \( S \). Then the flat section of the Levi-Civita connection of that metric defines the flat coordinate system (see e.g. \cite{44} for more details).

For weighted homogeneous polynomials, \( \{ \phi_{\alpha} d^{n}x \}_{\alpha} \subset \hat{H}_{F}^{(0)} \) is called a good basis if the vector subspace \( B = \text{Span}_{\mathbb{C}} \{ \phi_{\alpha} d^{n}x \}_{\alpha} \) satisfies \( K_{f}(B, B) \subset \mathbb{C}z^{n} \), where we note that the space \( B \) is isomorphic to the Jacobi algebra \( \text{Jac}(f) \) as \( \mathbb{C} \)-vector spaces. One key step to construct a primitive form is that the concept of the primitive forms is equivalent to the notion of good section \cite{42} (cf. \cite{46}). In order to show this, we need to extend a good basis in \( \hat{H}_{f}^{(0)} \) to a "deformed good basis" in the deformed module \( \hat{H}_{F}^{(0)} \), where, in the proof, we use a classical analytic result known as Birkhoff decomposition theorem. In \cite{28}, we replaced the role of the Birkhoff theorem by a multiplication of the "holomorphic part of the oscillatory integral factor" \( e^{f} \) : \( B \rightarrow B((z))[[s]] \) (here the first copy of \( B \) should be read off a subspace of the deformation \( \hat{H}_{F}^{(0)} \) of \( \hat{H}_{f}^{(0)} \)), which is able to calculate in power series in the local coordinate perturbatively. Inspired from this, we obtain the following, which is a combination of several propositions in section 3.2 of \cite{29}.
Proposition 2.1. Given a good basis \( \{ [\phi_{\alpha} dx^n] \}_{\alpha=1}^\mu \subset \mathcal{H}_f^{(0)} \), there exists a unique pair \((\zeta, \mathcal{J})\) satisfying the following: (1) \( \zeta \in \mathbb{B}[[z]][[s]], \) (2) \( \mathcal{J} \in [dx^n] + z^{-1} B[z^{-1}][[s]] \subset \mathcal{H}_f[[s]], \) and (\(\ast\))
\[
e^{(F - f)/z} \zeta = \mathcal{J}.
\]
Moreover, we embed \( z^{-1} \mathbb{C}[z^{-1}][[s]] \hookrightarrow z^{-1} \mathbb{C}[[z^{-1}]][[s]] \) and decompose
\[
\mathcal{J} = [dx^n] + \sum_{m=1}^{\infty} z^m \mathcal{J}_m, \quad \text{where} \quad \mathcal{J}_m = \sum_{\alpha} \mathcal{J}_m^\alpha [\phi_{\alpha} dx^n], \quad \mathcal{J}_m^\alpha \in \mathbb{C}[[s]].
\]
Then \( \zeta \) gives a formal primitive form, and \( \{ \mathcal{J}_m^\alpha \}_{m=2} \) give a formal Frobenius manifold structure on \( S \) with flat coordinates \( \{ \mathcal{J}_m^\alpha \}_m \). In particular, both \( \zeta \) and \( \mathcal{J} \) can be computed recursively by an algebraic algorithm via the above formula.

Explicitly, let us denote by \( J(\cdot, \cdot) \) and \( \ast \) the flat metric (the first residue pairing) and the product structure on the tangent bundle of \( S \), respectively. For simplicity, let us denote by \( t_1, \ldots, t_\mu \) the flat coordinate system on \( S \) and by \( \partial_{t_i}, \ldots, \partial_{t_\mu} \) their partial derivatives. Then, as a consequence of the flat structure, the following 3-tensor
\[
A(\partial_{t_i}, \partial_{t_j}, \partial_{t_k}) := J(\partial_{t_i} \ast \partial_{t_j}, \partial_{t_k}) = J(\partial_{t_i}, \partial_{t_j} \ast \partial_{t_k}) \in \Gamma(S, \mathcal{O}_S) \quad 1 \leq i, j, k \leq \mu
\]
is symmetric in the three variables, and satisfies the following integrability conditions
\[
\partial_{t_i} A(\partial_{t_j}, \partial_{t_l}, \partial_{t_k}) = \partial_{t_j} A(\partial_{t_l}, \partial_{t_i}, \partial_{t_k}) \quad \text{for all} \quad 1 \leq i, j, k, l \leq \mu.
\]
Therefore, there exists a function (formal power series in the flat coordinates) \( \mathcal{F}^{SG}_{0,f} \) on \( S \), called the prepotential, such that
\[
\partial_{t_i} \partial_{t_j} \partial_{t_k} \mathcal{F}^{SG}_{0,f} = A(\partial_{t_i}, \partial_{t_j}, \partial_{t_k}) = J(\partial_{t_i} \ast \partial_{t_j}, \partial_{t_k})
\]
(where the quadratic terms are normalized to be 0).

We are enabled to compute the prepotential \( \mathcal{F}^{SG}_{0,f} \) of the associated formal Frobenius manifold structure in a perturbative way, for an arbitrary weighted homogeneous singularity. On the other hand, it is shown in [25] that the formal power series \( \zeta \) is in fact the Taylor series expansion of the associated (analytic) primitive form around the origin \( 0 \in S \). This explains the geometric origin of the induced (formal) Frobenius manifold structure in the above proposition together with the analyticity of its prepotential \( \mathcal{F}^{SG}_{0,f} \).

Let us restrict our attention to the case of exceptional unimodular singularities now. Originally, there are 14 exceptional unimodular singularities by Arnold [11], which are one parameter families of singularities with three variables. Each family contains a weighted homogenous singularity characterized by the existence of only one negative degree but no zero-degree deformation parameter [43]. Hence in the present note, by exceptional
unimodular singularities, we mean the weighted homogeneous polynomials in these one parameter families, which are given in Table 1.

| Type | Polynomial | Polynomial | Polynomial | Polynomial |
|------|------------|------------|------------|------------|
| \(E_{12}\) | \(x^3 + y^7\) | \(W_{12}\) | \(x^4 + y^5\) | \(U_{12}\) | \(x^3 + y^3 + z^4\) |
| \(Q_{12}\) | \(x^2 y + x y^3 + z^3\) | \(Z_{12}\) | \(x^3 y + y^4 x\) | \(S_{12}\) | \(x^2 y + y^2 z + z^3 x\) |
| \(E_{14}\) | \(x^2 + x y^4 + z^3\) | \(E_{13}\) | \(x^3 + x y^5\) | \(Z_{13}\) | \(x^2 + x y^3 + y z^3\) | \(W_{13}\) | \(x^2 + x y^2 + y z^4\) |
| \(Q_{10}\) | \(x^2 y + y^4 + z^3\) | \(Z_{11}\) | \(x^3 y + y^5\) | \(Q_{11}\) | \(x^2 y + y^3 z + z^3\) | \(S_{11}\) | \(x^2 y + y^2 z + z^4\) |

There is a partial classification \[43\] of weighted homogeneous polynomial with isolated singularity by using the central charge

\[ \hat{c}_f := \sum_{i=1}^{n} (1 - 2q_i). \]

The case \(\hat{c}_f \leq 1\) is characterized as ADE-singularities if \(\hat{c}_f < 1\), or simple elliptic singularities if \(\hat{c}_f = 1\). The first examples of \(\hat{c}_f > 1\) are the exceptional unimodular singularities, the central charge of which are listed in Table 2 by direct calculations.

| Type | \(E_{12}\) | \(E_{13}\) | \(E_{14}\) | \(Z_{11}\) | \(Z_{12}\) | \(Z_{13}\) | \(W_{12}\) | \(W_{13}\) | \(Q_{10}\) | \(Q_{11}\) | \(Q_{12}\) | \(S_{11}\) | \(S_{12}\) | \(U_{12}\) |
|------|------------|------------|------------|---------|---------|---------|------------|------------|---------|---------|---------|---------|---------|--------|
| \(\hat{c}_f\) | 22/37 | 16/15 | 13/12 | 16/15 | 12/11 | 10/9 | 11/17 | 9/8 | 13/12 | 10/9 | 17/15 | 9/8 | 15/13 | 7/5 |

For the 14 singularities \(f\), the good basis is already known to be unique \[20, 28, 47\], and is simply given by a basis of Jacobian algebra \(\text{Jac}(f)\). Hence, the primitive form is unique (up to a nonzero scalar). By Proposition 2.1 we can obtain the data on LG B-model at genus zero in a perturbative way, and in particular we can calculate the four-point function \(\mathcal{F}_0^{(4)}\) (that is, the degree 4 terms of the prepotential \(\mathcal{F}_{0,f}^{SC}\) with respect to the flat coordinate system) of the Frobenius manifold structure associated to the primitive form. For instance for \(U_{12}\)-singularity, \(f = x^3 + y^3 + z^4\), we let \(\{\phi_i\}_i = \{1, z, x, y, z^2, xz, yz, xy, xz^2, yz^2, xyz, xyz^2\}\). By direct calculations, we obtain the four-point function in flat coordinates \((t_1, \ldots, t_{12})\) with respect to the primitive form \(\zeta = \ldots\)
\[ \mathcal{F}_0^{(4)} = \frac{1}{8}t_2^2t_6t_7 + \frac{1}{6}t_3^2t_8t_8 + \frac{1}{6}t_4^2t_7t_8 + \frac{1}{4}t_2^2t_5t_7t_9 + \frac{1}{6}t_3^2t_8t_9 + \frac{1}{4}t_2^2t_5t_6t_{10} + \frac{1}{6}t_4^2t_8t_{10} + \frac{1}{8}t_2^2t_9t_{10} + \frac{1}{8}t_2^2t_8t_{11} + \frac{1}{6}t_3^2t_6t_{11} + \frac{1}{6}t_4^2t_7t_{11} + \frac{1}{18}t_3^2t_{12} + \frac{1}{18}t_4^2t_{12} + \frac{1}{18}t_2^2t_8t_{12} \]

Here we make boxes for the last three monomials, which will be compared with the data on the LG \(A\)-side, studied in the next section.

3. MIRROR CONSTRUCTION AND THE FAN-JARVIS-RUAN-WITTEN THEORY

For the mirror symmetry purpose, we restrict our singularity into an invertible polynomial, where the number of variables is the same as the number of monomials in the polynomial. We consider a pair \((W, G)\), where \(W\) is an invertible polynomial with \(n\) variables \(x_1, \cdots, x_n\) and has no monomials of the form \(x_i x_j\) for \(i \neq j\). By rescaling the variables, we can always write this polynomial by

\[ W = \sum_{i=1}^{n} \prod_{j=1}^{n} x_i^{a_{ij}}. \]

The matrix \(E_W := (a_{ij})_{n \times n}\) of exponents is called the exponent matrix of \(W\). Let us use \(\text{Aut}(W)\) to denote the group of diagonal symmetries of \(W\),

\[ \text{Aut}(W) := \{ \text{diag}(\lambda_1, \cdots, \lambda_n) \mid W(\lambda_1 x_1, \cdots, \lambda_n x_n) = W(x_1, \cdots, x_n), \lambda_i \in \mathbb{C}^* \} . \]

Then \(G\) is a subgroup in \(\text{Aut}(W)\) containing

\[ J_W := \text{diag} \left( \exp(2\pi \sqrt{-1} q_1), \cdots, \exp(2\pi \sqrt{-1} q_n) \right), \]

with \(q_1, \cdots, q_n\) are the weights of variables in \(W\). Berglund and Hübsch constructed a mirror polynomial \(W^T\) \([2]\) by taking

\[ W^T = \sum_{i=1}^{n} \prod_{j=1}^{n} x_j^{a_{ji}} \]

where \(E_W^T\) is the transpose matrix of \(E_W\). In general, the LG-LG mirror symmetry relates the pair \((W, G)\) to a mirror pair \((W^T, G^T)\), where \(G^T\) is constructed by \([3, 20]\).

In particular, if \(G = \text{Aut}(W)\), then \(G^T\) is the group with only an identity element. A LG-LG mirror symmetry conjecture can be formulated as the equivalence of Frobenius manifold structure associated with the primitive form theory of \(W^T\) and that associated with the genus-0 Fan-Jarvis-Ruan-Witten theory (FJRW) theory of \((W, G = \text{Aut}(W))\).

The FJRW theory is introduced by Fan, Jarvis and Ruan in a series of papers \([14, 15]\), based on a proposal of Witten \([52]\). The theory works for the pair \((W, G)\) in general, where \(W\) is a weighted homogenous polynomial which has an isolated critical point at
the origin and $G$ is a subgroup in $\text{Aut}(W)$. The theory also requires that $G$ contains $J_W$. For technical reasons, $W$ does not contain any monomial term $xy$. In the present note, we will focus only on the case $G = \text{Aut}(W)$.

For a pair $(W, \text{Aut}(W))$, there is an FJRW state space $H_W$ which collects all $\text{Aut}(W)$-invariant part of middle dimensional Lefschetz thimble on the fixed locus of each group element $\gamma$ in $\text{Aut}(W)$,

\[ H_W := \bigoplus_{\gamma \in \text{Aut}(W)} H^{\text{mid}}(\text{Fix}(\gamma); W^\infty, \mathbb{C})^{\text{Aut}(W)}. \]

Here $W^\infty$ is the preimage of $[M, \infty)$, for $M \gg 0$, under the real part of $W$ restricted on the fixed locus $\text{Fix}(\gamma)$.

Fan, Jarvis and Ruan [14, 15] studied the space of solutions of Witten equations for

\[ \frac{\partial u_i}{\partial z} + \frac{\partial_i W(u_1, \cdots, u_n)}{u_i} = 0, \quad i = 1, \ldots, n \]

where $z$ is a local coordinate of the curve in consideration (but not the formal variable in primitive form theory) and $u_i$ ($1 \leq i \leq n$) is a section of a line bundle $L_i$ with suitable degrees over the curve (for algebraic construction, see [5, 34]), and constructed a cohomological field theory (in the sense of Kontsevich-Manin [25])

\[ \{ \Lambda_{g,k}^W : (H_W)^{\otimes k} \to H^*(\overline{M}_{g,k}, \mathbb{C}) \} \]

on moduli space of stable curves $\overline{M}_{g,k}$. As a consequence, this gives the FJRW invariants

(3.1) \[ \langle \alpha_1 \psi_1^{(i_1)}, \ldots, \alpha_k \psi_k^{(i_k)} \rangle_{g,k}^W = \int_{\overline{M}_{g,k}} \Lambda_{g,k}^W(\alpha_1, \ldots, \alpha_k) \prod_{j=1}^{k} \psi_j^{(i_j)}, \quad \alpha_j \in H_W. \]

Here $\psi_j$ is the $j$-th psi-class on $\overline{M}_{g,k}$. The genus-0 invariants without $\psi$-class involved give a formal Frobenius manifold structure on $H_W$. The prepotential of this formal Frobenius manifold is

\[ F_{FJRW}^{0} = \sum_{k \geq 3} \frac{1}{k!} (t_0, \ldots, t_k)^{W}_{0,k}, \quad t_0 = \sum_{j=1}^{\mu} t_{0,\alpha_j} \alpha_j. \]

It is a formal power series of $t_{0,\alpha_j}, j = 1, \cdots, \mu$. More generally, the FJRW total ancestor potential $\mathcal{A}^{FJRW}_W$ is defined to be

\[ \mathcal{A}^{FJRW}_W = \exp \left( \sum_{k \geq 0} \frac{1}{k!} \langle t(\psi_1) + \psi_1, \ldots, t(\psi_k) + \psi_k \rangle_{g,k}^W \right). \]

Here $t(z) = \sum_{m \geq 0} \sum_{j=1}^{\mu} t_{m,\alpha_j} \alpha_j z^m$.

4. MIRROR SYMMETRY FOR EXCEPTIONAL UNIMODULAR SINGULARITIES

In [29], the following isomorphism between two types of Frobenius manifolds is proven.
**Theorem 4.1.** Let $W^T$ be one of the 14 exceptional unimodular singularities in Table 1. There exists a mirror map \( \Psi : \text{Jac}(W^T) \cong H_W \), which induces an equality

\[
F_{0,W}^{SG} = F_{0,W}^{FJRW}.
\]

The mirror map \( \Psi : \text{Jac}(W^T) \to H_W \) is constructed by Krawitz \[26\] and proven that it is a ring isomorphism under a technical condition that \( W \) (in the FJRW side) is not allowed to be a chain type polynomial with one weight \( 1/2 \). For exceptional unimodular singularities, this condition excludes two examples, \( W^T = x^2y + y^3z + z^3(Q_{11}) \), \( x^2y + y^2z + z^4(S_{11}) \). However, in \[29\], the technical condition is removed by using the Getzler’s relation in \( M_{1,4} \) \[17\].

The proof of Theorem 4.1 mainly uses the axioms of cohomological field theories, in particular, the Witten-Dijkgraaf-Verlinde-Verlinde (WDVV) equations. Combined with the special properties on the weights of the exceptional unimodular singularities, it was proved in \[29\] that both \( F_{0,W}^{SG} \) and \( F_{0,W}^{FJRW} \) are determined by the underlying ring isomorphism and a few initial invariants \( ⟨· · ·⟩_{0,4} \). The invariants on the primitive form theory side can be calculated by the perturbative formula. On the other hand, again by some WDVV equations, the invariants on the FJRW side can be reduced to invariants which can be calculated by the so-called orbifold-Grothendieck-Riemann-Roch formula. Under the mirror map which identifies the deformation parameter space in primitive form side to the FJRW state space together with the ring structure and the inner product, the invariants on both sides are identified up to a scale \( -1 \). Then, by rescaling mirror map appropriately, we obtain the desired equality (4.1).

This equality of the pre-potentials in genus 0 is lifted to the equality of higher genus potentials as follows. For the generic point \( s \in S \) in the universal unfolding \( F \) of \( W^T \), the \( F(x,s) \) is a Morse function in \( x \) so that its Jacobian ring is a direct sum of the one dimensional algebra \( \mathbb{C} \). That is, after the Kodaira-Spencer map identification, the Frobenius algebra structure on the tangent space of \( S \) at \( s \) is *semi-simple*. Such a generic point is called *semisimple*. Givental defined a *total ancestor potential* (or a higher genus formula) \[18\] using only the genus zero data near the generic point together with the knowledge of the Witten-Kontsevich tau-function. The later is just also called the *total ancestor potential* of the Gromov-Witten theory with the target being a point. Teleman \[51\] proved that this higher genus formula in a cohomological field theory is uniquely determined by the underlying Frobenius manifold at the semisimple point. The origin in the universal unfolding space \( S \) is not semisimple, however, Givental’s formula can be uniquely extended to \( σ_{W}^{SG} \) at the origin by Milanov \[31\] (see also Coates-Iritani \[9\]). The uniqueness of the
extension will upgrade Theorem 4.1 to an identity of higher genus potential function:

$$\mathcal{A}_{WG}^{SG} = \mathcal{A}_{W}^{FJRW}.$$ 

This completes a proof of LG-LG mirror symmetry.

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