The dust energy balance in the edge-on spiral galaxy NGC 4565

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ABSTRACT

We combine new dust continuum observations of the edge-on spiral galaxy NGC 4565 in all Herschel/Spectral and Photometric Imaging Receiver (250, 350 and 500 µm) wavebands, obtained as part of the Herschel Reference Survey, and a large set of ancillary data (Spitzer, Sloan Digital Sky Survey, Galaxy Evolution Explorer) to analyse its dust energy balance. We fit a radiative transfer model for the stars and dust to the optical maps with the fitting algorithm FITSKIRT. To account for the observed ultraviolet and mid-infrared emission, this initial model was supplemented with both obscured and unobscured star-forming regions. Even though these star-forming complexes provide an additional heating source for the dust, the far-infrared/submillimetre emission long wards of 100 µm is underestimated by a factor of 3–4. This inconsistency in the dust energy budget of NGC 4565 suggests that a sizable fraction (two-thirds) of the total dust reservoir ($M_d \sim 2.9 \times 10^8 M_\odot$) consists of a clumpy distribution with no associated young stellar sources. The distribution of those dense dust clouds would be in such a way that they remain unresolved in current far-infrared/submillimetre observations and hardly contribute to the attenuation at optical wavelengths. More than two-thirds of the dust heating in NGC 4565 is powered by the old stellar population, with localized embedded sources supplying the remaining dust heating in NGC 4565. The results from this detailed dust energy balance study in NGC 4565 are consistent with that of similar analyses of other edge-on spirals.

Key words: radiative transfer – dust, extinction – galaxies: individual: NGC 4565 – galaxies: ISM – infrared: galaxies.

1 INTRODUCTION

Panchromatic radiative transfer 2D and 3D modelling of galaxies provide a powerful tool to analyse the characteristics of dust in galaxies (i.e. optical properties, distribution, clumpiness, etc.) in a self-consistent way (e.g. Xilouris et al. 1999; Popescu et al. 2000, 2011; Alton et al. 2004; Bianchi 2008; Baes et al. 2010; MacLachlan et al. 2011; De Looze et al. 2012; Holwerda et al. 2012). From optical observations, the properties and spatial distribution of stars and dust can be constrained using a radiative transfer code to model the
propagation of stellar light and its interaction with dust particles in a galaxy. In a second step, the dust emission predicted from the radiative transfer simulations is compared to the observed thermal dust re-emission at infrared (IR)/submillimetre (submm) wavelengths. Such a complementary study imposes that dust features can easily be identified from optical as well as IR observations. This requirement has limited the number of galaxies for which detailed dust energy balance studies have been attempted in the past. Edge-on spirals are considered ideal cases for those complementary studies since projection effects allow to resolve the dust distribution vertically (e.g. Misiriotis et al. 2001; Alton et al. 2004; Dasyra et al. 2005; Bianchi 2008; Baes et al. 2010; Popescu et al. 2011; De Looze et al. 2012). Unfortunately, the large inclination angle impedes the characterization of substructures in the dust organization and the identification of heavily obscured star-forming regions. Only the application of complex 3D models in conjunction with competitive fitting algorithms have been shown to successfully characterize clumpiness and asymmetries in the stellar and dust distribution of highly inclined objects (e.g. Schechtman-Rook, Bershady & Wood 2012).

Dust energy balance studies of individual edge-on galaxies reveal an inconsistency between the predicted far-infrared (FIR)/submm fluxes of radiative transfer models and the observed emission in those wavebands (e.g. Popescu et al. 2000, 2011; Misiriotis et al. 2001; Alton et al. 2004; Dasyra et al. 2005; Baes et al. 2010; De Looze et al. 2012; Holwerda et al. 2012). Although radiative transfer models might successfully explain the observed optical attenuation, the modelled dust emission underestimates the observed thermal dust re-emission by a factor of 3–4. In order to reconcile the results of the radiative transfer models with the observations, two scenarios have been proposed: either a significant underestimation of the FIR/submm dust emissivity has been argued (e.g. Alton et al. 2004; Dasyra et al. 2005; MacLachlan et al. 2011) or, alternatively, the distribution of a sizable fraction of dust in clumps or a second inner dust disc having a negligible attenuation on the bulk of the starlight (e.g. Popescu et al. 2000; Misiriotis et al. 2001; Bianchi 2008; MacLachlan et al. 2011; De Looze et al. 2012). However, a general consensus on the origin of the energy balance problem has not yet been achieved.

Until recently, observations of dust in galaxies were hampered by the poor resolution and limited wavelength coverage of IR instrumentation. The first FIR space satellites [e.g. IRAS, Infrared Space Observatory (ISO), Spitzer, Akari] only covered mid-infrared (MIR) and FIR wavelengths up to 200 µm. Most ground-based submm/mm instrumentation either only start operating from 850 µm onwards [e.g. the Large APEX Bolometer Camera (LABOCA), IRAM] or suffer severely from the earth’s opacity at shorter submm wavelengths [450 µm; e.g. Submillimetre APEX Bolometer Camera (SABOCA), Submillimeter Common User Bolometer Array (SCUBA)]. This coarse sampling of the Rayleigh–Jeans tail of dust emission in most galaxies induced poor constraints on the temperature, content and properties of dust in those objects. Recently, this large void in wavelength coverage was bridged thanks to the launch of the Herschel Space Observatory (Pilbratt et al. 2010), offering the opportunity to trace the dust emission from galaxies in a broad wavelength domain ranging from 55 to 672 µm. The high spatial resolution achieved by Herschel is providing a large collection of galaxies suitable for detailed dust energy balance analyses, supplemented with the ongoing development of powerful radiative transfer codes.

At a distance of about 16.9 Mpc, the Needle Galaxy (NGC 4565) is one of the most nearby edge-on spiral galaxies covering an angular scale of ~15 arcmin on the sky. Distance measurements for NGC 4565 range from 9 to 22 Mpc depending on the applied technique. We adopt a value of $D \sim 16.9$ Mpc throughout this paper, which is the average of the most recent distance measurements from $J$-band surface brightness fluctuations as reported in Tonry et al. (2001) and Jensen et al. (2003). All values from the literature referred to in this paper were converted to the assumed distance of 16.9 Mpc. NGC 4565 is a Sb spiral, classified as a Seyfert galaxy (Ho et al. 1997) and located in the Coma I cloud. Recently its classification as an active galactic nucleus (AGN) has been confirmed in the MIR with the detection of the [Ne v] lines with Spitzer (Laine et al. 2010), which are among the strongest indicators of AGN (see e.g. Tommasi et al. 2010). The stellar geometry in NGC 4565 has been studied extensively at optical and near-infrared (NIR) wavelengths (van der Kruit 1979; Hamabe et al. 1980; van der Kruit & Searle 1981; Jensen & Thuan 1982; Dettmar & Wielebinski 1986; Rice et al. 1996; Naeslund & Jørsaaeter 1997; Wu et al. 2002; Seth, Dalcanton & de Jong 2005; Kormendy & Barentine 2010; Mosenkov, Sotnikova & Reshetnikov 2010). From studies of the absorption layer in NGC 4565, Ohta & Kodaira (1995) claim a dust distribution in a ring-like structure and an inner cut-off radius along the major axis at ~130 arcsec. At radii smaller than the cut-off radius dust obscuration effects seem to be absent, either due to a deficiency of dust or the distribution of a central dust component preventing the identification of dust attenuation from optical data (Ohta & Kodaira 1995). Studies of residual J- and K-band images obtained after fitting the stellar component and subtracting those models from the NIR images in Rice et al. (1996) result in estimates for the dust scale length ~66 arcsec and scale height ~5 arcsec in NGC 4565, which corresponds to about 60 and 50 per cent of the stellar disc in the K band.

In addition to these optical/NIR studies of the dust component, the properties and geometry of the dust in NGC 4565 have been analysed as well from IR/submm observations (Wainscoat, de Jong & Wesselius 1987; Engargiola & Harper 1992; Neininger et al. 1996; Rice et al. 1996; Alton et al. 2004; Kormendy & Barentine 2010; Laine et al. 2010). In the MIR Infrared Array Camera (IRAC) 8-µm waveband, Kormendy & Barentine (2010) and Laine et al. (2010) identify a ring of prominent polycyclic aromatic hydrocarbon emission from dust at the same position as the molecular gas ring at radii of ~80–100 arcsec (Sofue & Nakai 1994; Neininger et al. 1996). The first FIR study on dust emission in NGC 4565 based on IRAS data calculate a total IR luminosity $L_{\text{IRas}} \sim 3 \times 10^9 L_\odot$ (Wainscoat et al. 1987), which compares well to the dust emission from our Galaxy. The dust emission in NGC 4565 however is more extended compared to the dust distribution in our Galaxy (Wainscoat et al. 1987). Based on estimates of the IR excess, Wainscoat et al. (1987) argue that only one quarter of the dust reservoir in NGC 4565 is heated by embedded stars in molecular clouds with the remainder of the dust heating power provided by the diffuse interstellar radiation field.

Observations at FIR wavelengths (100, 160 and 200 µm) with the Kuiper Airborne Observatory (KAO; Engargiola & Harper 1992) find a two-component dust model with a cold dust reservoir ($T_d \sim 20$ K) distributed in an exponential disc and a warmer dust component ($T_d \sim 30$ K) spatially coinciding with a bisymmetric spiral pattern in NGC 4565. In spatial correlation with the more extended ring of H I gas, Neininger et al. (1996) also found a colder dust component (~15 K) in the outskirts of the galaxy. A second dust reservoir ($T_d \sim 18$ K) was found in the centre of NGC 4565 and in the molecular gas ring at radii of ~80–100 arcsec. The central concentration of both dust and molecular gas (Neininger et al. 1996)
contradicts the suggested absence of dust in the centre of the galaxy (Ohta & Kodaira 1995).

Alton et al. (2004) constrain the amount of dust in NGC 4565 from a complementary study of the attenuation properties of dust in the optical and the thermal dust re-emission at 1.2 mm. From optical constraints, they construct a radiative transfer model for the stars and dust in NGC 4565. Upon comparison of the visual optical depth in their radiative transfer model with thermal continuum radiation in the 1.2-mm waveband, Alton et al. (2004) infer a dust emissivity at 1.2 mm which is 1.5 times higher than the benchmark, semi-empirical model from Draine & Lee (1984) used in the radiative transfer simulation. From their combined study on the dust emissivity in NGC 4565 and two other nearby spirals, they argue that coagulation of well-ordered dust crystalline grains into amorphous particles in high-density environments can influence the emissivity of dust at submm/mm wavelengths.

The first submm observations of NGC 4565 were obtained during the first test flight of the 2 m Balloon-borne Large Aperture Submillimeter Telescope (BLAST) in 2005 (Wiebe et al. 2009), covering the galaxy in the same wavebands (250, 350 and 500 µm) as offered by the Spectral and Photometric Imaging Receiver (SPIRE) instrument onboard the Herschel Space Observatory. Although the resolution in their final images was degraded by the large point spread functions (PSFs) (∼186–189 arcsec; Truch et al. 2008), Wiebe et al. (2009) could estimate the scale length of the dust lane hR ∼ 150 arcsec from the beam-convolved image at 250 µm, which is about twice the value obtained from optical studies (Rice et al. 1996). A modified blackbody fit to the BLAST fluxes and additional constraints from ISO (170 µm; Stickel et al. 2004) and IRAS (12, 25, 60 and 100 µm; Sanders et al. 2003) observations predict a dust reservoir of mass Md ∼ 3 × 10^8 M⊙ (scaled to a distance of 16.9 Mpc) at a temperature Td ∼ 16 K for an emissivity index β = 2 and a Draine & Li (2007) dust grain composition (Wiebe et al. 2009).

NGC 4565 was recently observed in all SPIRE wavebands (250, 350 and 500 µm) as part of the Herschel Reference Survey (HRS; Boselli et al. 2010). In this work, we will analyse the dust characterisation from a spatially resolved study of the Herschel IR/submm photometry and a wealth of ancillary observations. Section 2 reviews the observing and data reduction strategy for the Herschel observations and the set of ancillary data used in this study. Section 3 presents a detailed study of the dust energy balance in NGC 4565 with a brief introduction to the applied radiative transfer code SKIRT and fitting algorithm FITSKIRT (Section 3.1) and a gradual construction of a self-consistent model that can account for the observed quantities for NGC 4565 across the multiwavelength spectrum (Section 3.2). In Section 4, the results of our radiative transfer modelling procedure are discussed and compared to the dust energy balance, gas-to-dust ratio and dust heating mechanisms in other nearby galaxies.

2 OBSERVATIONS AND DATA REDUCTION

2.1 Herschel data (250–500 µm)

NGC 4565 was observed on the 2009 December 31 (ObsID 0 × 500002cd8) as part of the HRS (Boselli et al. 2010), a Herschel Guaranteed Time Key Program observing 323 galaxies in the nearby Universe spanning a wide range in morphological type and environment. Dust continuum maps were obtained in all SPIRE wavebands (centred at 250, 350 and 500 µm), covering the galaxy in three cross-linked scans (nominal and orthogonal) at medium scan speed (30 arcsec s⁻¹). The final maps cover a square area of 30 × 30 arcmin², corresponding to an extent of at least 1.5 times the optical isophotal diameter D25 in NGC 4565.

The SPIRE data were processed up to Level-1 in the Herschel Interactive Processing Environment (HIPE, version 4.0.1367: Ott 2010) with the standard script adapted from the official pipeline (POD5 pipeline.py, dated 2010 June 8) as provided by the SPIRE Instrument Control Centre (ICC). Instead of the ICC default settings for deglitching, we used the WAVELETDEGLITCHER, which was adjusted to mask the sample following a glitch. The WAVELETDEGLITCHER algorithm was applied a second time after the flux calibration, since this was shown to significantly improve the removal of remaining glitches for the version of the code at that time (currently it is only run once). Instead of the pipeline default temperature drift correction and median baseline subtraction, we applied a custom method (BriGAdE; Smith et al., in preparation) to remove the temperature drift and bring all bolometers to the same level. Final SPIRE maps (see Fig. 1, last three bottom panels) were created with the naive mapper in HIPE with pixel sizes of 6, 8 and 12 arcsec at 250, 350 and 500 µm, respectively. The full width at half-maximum (FWHM) of the SPIRE beams have sizes of 18.2, 24.5 and 36.0 arcsec at 250, 350 and 500 µm (for pixel sizes of 6, 8 and 12 arcsec, respectively. To update our flux densities to the latest v8 calibration product, the 350 µm images are multiplied by a factor of 1.0067 (SPIRE Observers’ Manual 2011). Since the calibration in the standard data reduction pipeline was optimized for point sources, correction factors (0.9828, 0.9834 and 0.9710 at 250, 350 and 500 µm) are applied to convert the K4 colour correction factors from point source to extended source calibration (SPIRE Observers’ Manual 2011). Appropriate colour correction factors were determined from a single component modified blackbody fit to the Rayleigh–Jeans tail of dust emission [constrained by fluxes from Multiband Imaging Photometer for Spitzer (MIPS) 160 µm and all SPIRE wavebands] for values β = 1.0, 1.5 and 2.0 (i.e. the range of spectral indices derived for the HRS sample galaxies; Boselli et al. 2012). With the best fit obtained for values β = 2 and Td = 17 K, we apply colour correction factors (1.019, 1.009 and 1.021) to the 250, 350 and 500 µm images (SPIRE Observers’ Manual 2011), respectively, to correct for the shape of the response function in every filter. An uncertainty in the calibration of 7 per cent was assumed (Swinyard et al. 2010; SPIRE Observers’ Manual 2011).

From aperture photometry within an elliptical aperture (20 × 2.7 arcmin²) flux densities of 62.06, 31.07 and 12.87 Jy for NGC 4565 were obtained at 250, 350 and 500 µm, respectively (Ciesla et al. 2012). Since those fluxes were corrected for extended source calibration, we only need to correct those values with appropriate colour corrections factors (1.019, 1.009 and 1.021) to the 250, 350 and 500 µm images (SPIRE Observers’ Manual 2011), respectively. The large ancillary data set of FIR/submm observations allows us to compare those SPIRE fluxes to flux measurements from other space or airborne missions in overlapping wavebands. Wiebe et al. (2009) report fluxes obtained with BLAST at 250 µm (37.2 ± 4.5 Jy), 350 µm (21.0 ± 2.1 Jy) and 500 µm (9.8 ± 0.9 Jy), which are 40, 32 and 24 per cent lower with respect to SPIRE flux measurements (see Table 1). BLAST observations of NGC 4565 covered an area of ~0.4 deg² centred on the galaxy, sufficient to detect the dust emission from NGC 4565 in those wavebands. Considering that BLAST fluxes for NGC 4565 were measured from data taken during the first BLAST test flight in 2005,
Figure 1. The stellar and dust emission from NGC 4565 as observed from GALEX FUV and NUV, SDSS ugriz, Spitzer IRAC 3.6, 5.8 and 8.0 µm, Spitzer MIPS 24, 70 and 160 µm and Herschel SPIRE 250, 350 and 500 µm images. The IRAC 4.5-µm map is not explicitly shown here because of the great resemblance with the IRAC 3.6-µm image. The beam sizes of all instruments in the IR/submm wavebands have been indicated as white circles in the bottom left-hand corner of the corresponding panels.

the images were degraded by the large PSF of BLAST05 (FWHM of ~186–189 arcsec; Truch et al. 2008). Prior to aperture photometry, the images were deconvolved following the method outlined in Chapin et al. (2008). This beam deconvolution might be the cause of the lower BLAST fluxes compared to SPIRE measurements. Alternatively, calibration issues might be responsible for the small offset between BLAST and SPIRE. Also Planck measurements (Planck Collaboration et al. 2011) at 350 µm (24.35 ± 1.43 Jy) and 550 µm (8.12 ± 0.48 Jy) are about 20 per cent below our SPIRE measurements. The lower Planck values can easily be explained from the radius applied for aperture photometry on point sources in their first release catalogue (Planck Collaboration et al. 2011). Indeed, an aperture with radius fixed to the FWHM of Planck at 350 µm (~4.23 arcmin) does not cover the entire galaxy’s emission in that waveband (see the upper right-hand panel of fig. 10 in Ciesla et al. 2012). Performing aperture photometry within the Planck radius at 350 µm on the SPIRE images results in a flux density $S_\nu \sim 25.5 \pm 0.9$ Jy (Ciesla et al. 2012), which is in close agreement with the reported Planck value (Planck Collaboration et al. 2011).

Fig. 1 (last three bottom panels) shows the Herschel maps for NGC 4565 in the SPIRE 250, 350 and 500 µm wavebands. Weak warping signatures at the edges of the dust disc can be identified in all Herschel dust maps, with the warp on the north-western side of the galaxy’s disc being more pronounced. These warp features
Aside from the warps on the edges of the galaxy, three peaks in the dust emission profile along the major axis from 24-μm and 1.2 mm dust continuum observations, smoothed to a common resolution of 20 arcsec. Based on these major axis profiles, we can identify a depression up to the 80-arcsec radius from the centre in the H\textsc{i} profile and a more or less constant column density profile in the outskirts of the galaxy from that radius onwards. The CO profile in NGC 4565 indicates the presence of a large molecular gas reservoir in the inner regions, distributed in a central disc and narrow ring with peak density around 80 arcsec, where also the inner depression in H\textsc{i} density ends. The SPIRE 250-μm dust emission profile along the major axis correlates well with the molecular gas density profile (see also Wainscoat et al. 1987; Rice et al. 1996; Kormendy & Barentine 2010; Laine et al. 2010) as well as the column density of the combined H\textsc{i} and H$_2$ gaseous components. The latter column density profile seems dominated by the peaks in the molecular gas component, rather than the smooth H\textsc{i} distribution across the plane of the galaxy. In contrast to the CO profile reported in Neininger et al. (1996), the CO observations from Sofue & Nakai (1994) identified a more asymmetric CO distribution with the majority of the molecular gas reservoir residing on the north-west side of the galaxy. The dust emission from the colder dust reservoir in the SPIRE 250-μm waveband however does not show a tendency for asymmetry.

### Table 1. Panchromatic overview of flux densities used in the dust energy balance study of NGC 4565.

| Filter | λ (μm) | $F_J$ (Jy) | Ref$^a$ |
|-------|--------|------------|--------|
| FUV   | 0.15   | 0.008 ± 0.001 | 1 |
| NUV   | 0.23   | 0.015 ± 0.002 | 1 |
| SDSS u| 0.36   | 0.07 ± 0.01  | 1 |
| SDSS g| 0.47   | 0.33 ± 0.03  | 1 |
| SDSS r| 0.62   | 0.70 ± 0.07  | 1 |
| SDSS i| 0.75   | 1.08 ± 0.10  | 1 |
| SDSS z| 0.89   | 1.45 ± 0.14  | 1 |
| 2MASS J| 1.25  | 2.18 ± 0.03  | 3 |
| 2MASS H| 1.65  | 3.05 ± 0.25  | 3 |
| 2MASS K| 2.20  | 5.01 ± 0.44  | 3 |
| WISE 1| 3.4    | 1.54 ± 0.04  | 3 |
| IRAC 3.6| 3.6  | 1.39 ± 0.14  | 5 |
| IRAC 4.5| 4.5  | 0.91 ± 0.09  | 5 |
| IRAC 5.8| 5.8  | 0.90 ± 0.03  | 4 |
| IRAC 8.0| 8.0  | 1.17 ± 0.14  | 5 |
| IRAC 12.1| 12.1 | 2.20 ± 0.05  | 4 |
| IRAS 12| 12    | 1.53 ± 0.23  | 6 |
| IRAS 22.2| 22.2 | 1.60 ± 0.16  | 4 |
| MIPS 24| 24    | 1.65 ± 0.07  | 6 |
| IRAS 25| 25    | 1.70 ± 0.26  | 6 |
| IRAS 60| 60    | 9.83 ± 1.47  | 6 |
| MIPS 70| 70    | 19.3 ± 1.9   | 6 |
| IRAS 100| 100  | 47.23 ± 7.08 | 6 |
| KUIPER 100| 100 | 52.00 ± 10.40 | 8 |
| MIPS 160| 160  | 86 ± 10      | 7 |
| KUIPER 160| 160 | 62.00 ± 12.40 | 8 |
| MIPS 200| 200   | 55.00 ± 11.00 | 8 |
| SPIRE 250| 250  | 63.24 ± 1.18 | 9 |
| SPIRE 350| 350  | 31.35 ± 0.69 | 9 |
| SPIRE 500| 500  | 13.14 ± 0.29 | 9 |
| IRAM 1200| 1200| 1.00 ± 0.10  | 10 |

$^a$References: 1 – Cortese et al. (2012b); 2 – Cortese (private communication); 3 – Jarrett et al. (2003); 4 – this paper, see Section 2.3; 5 – Ciesla et al. (in preparation); 6 – Rice et al. (1988); 7 – Bendo, Galliano & Madden (2012b); 8 – Engargiola & Harper (1992); 9 – Ciesla et al. (2012); 10 – Neininger et al. (1996).

Minor discrepancies in the column density profiles might also occur due to current limitations in the CO observations. Asymmetry in the dust emission peaks were identified in NGC 4565, with the CO observations from Sofue & Nakai (1994) and the CO observations from Sofue & Nakai (1994) identified a more asymmetric CO distribution with the majority of the molecular gas reservoir residing on the north-west side of the galaxy. The dust emission from the colder dust reservoir in the SPIRE 250-μm waveband however does not show a tendency for asymmetry.

### 2.2 Spitzer data (3.6–160 μm)

NGC 4565 was observed with the IRAC and MIPS instruments onboard the Spitzer Space Observatory as part of the programs Brown dwarf galaxy haloes (PI: Giovanni Fazio; AOR key: 3627776) and The formation of dust lanes in nearby, edge-on disc galaxies (PI: Roelof S. de Jong; AOR key: 14481408), respectively. IRAC (3.6, 4.5, 5.8 and 8.0 μm) and MIPS (24, 70 and 160 μm) data were retrieved from the Spitzer archive and reduced following the procedure outlined in Ciesla et al. (in preparation) and Bendo et al. (2012b). Although Spitzer IRAC data of NGC 4565 were already used in the analyses of Kormendy & Barentine (2010) and Laine et al. (2010), the maps presented in this work were processed independently. Final IRAC images have pixel sizes of 0.6 arcsec, whereas MIPS images were processed to final maps with pixel sizes of 1.5, 4.5 and 9.0 arcsec at 24, 70 and 160 μm, respectively. The IRAC beam has an almost uniform size in all wavebands (mean FWHM ~ 1.7, 1.7, 1.9 and 2.0 arcsec at 3.6, 4.5, 5.8 and 8 μm; Fazio et al. 2004). The FWHM of the MIPS beam varies from 6 arcsec at 24 μm, 18 arcsec at 70 μm to 38 arcsec at 160 μm (Engelbracht et al. 2007; Gordon et al. 2007; Stansberry et al. 2007).
wavebands (Reach et al. 2005). For MIPS, calibration uncertainties of 4, 10 and 12 per cent were reported for the 24, 70 and 160 µm wavebands (Engelbracht et al. 2007; Gordon et al. 2007; Stansberry et al. 2007).

MIPS fluxes at 24 and 70 µm compare reasonably well with IRAS observations of NGC 4565 at 25 and 60 µm (see Table 1). The longer wavelength MIPS observations at 160 µm are furthermore consistent within 30 per cent with data from the KAO in the same waveband (see Table 1).

Fig. 1 (last panel on the second row) shows the stellar emission in the IRAC 3.6 µm image. Dust emission from IRAC (5.8 and 8 µm) and MIPS observations are presented in Fig. 1 (third row and first panel on fourth row). The emission in the IRAC 3.6 and 4.5 µm wavebands is dominated by the stellar emission from the bulge in NGC 4565. The stellar emission from the bulge is also discernible in the IRAC 5.8 µm image. At 5.8 and 8 µm, emission from polycyclic aromatic hydrocarbons (PAHs) and very small grains (VSGs) distributed in a ring centred at 80–100 arcsec dominates. The presence of a hot dust component heated predominantly by star formation is also clearly detected in the MIPS 24-µm image. Besides this ring of hot dust, the central region of NGC 4565 is also a prominent emission source in the IRAC 5.8, 8 µm and MIPS 24 µm wavebands.

This central dust concentration in NGC 4565 is most likely heated by the AGN, since the inner 80-arcsec region of NGC 4565 seems to lack any star formation activity (see Section 2.4). Owing to the coarser resolution at 70 and 160 µm, the dust emission features are not as easily identified compared to shorter wavelength data. To better contrast the structures in the FIR dust emission, we overlaid the contours of dust emission on those FIR MIPS data. The contours show three peaks in the FIR dust emission of NGC 4565, similar to the peaks identified from the Herschel data, with prominent emission from a ring at 80–100 arcsec and a central dust concentration. Also the dust warps are manifested on the edges of the dust disc, in particular on the north-western side, in both IRAC (5.8, 8 µm) and all MIPS wavebands.

In a similar way as for the SPIRE observations at 250 µm, the resolution in the MIPS 24-µm map was degraded to 20 arcsec, allowing an immediate comparison with the distribution of atomic and molecular gas along the major axis of NGC 4565 (see Fig. 2). With the shape of the PSF in the MIPS 24-µm filter deviating from a Gaussian distribution, appropriate kernels were constructed following the procedure outlined in Bendo et al. (2012a) to degrade the resolution in the MIPS 24-µm image. With the 24 µm emission tracing the warm dust component in NGC 4565, we find a spatial correspondence with the dust heated by young stellar objects in NGC 4565 and the molecular gas reservoir, which provides the birth material for this star formation. The resemblance between the MIPS 24 µm, SPIRE 250 µm and IRAM 1.2 mm dust profiles suggests that the cold and warm dust component are well mixed in the interstellar medium (ISM) of NGC 4565. However, the asymmetry in CO observations from Sofue & Nakai (1994) does become evident in the dust emission profiles of the warmer dust component. This asymmetry is most likely the consequence of spiral arm structures in the plane of the galaxy. Similar asymmetric absorption features could also be identified in optical images (Hamabe et al. 1980; Rice et al. 1996) and were assigned to a spiral arm structure at position angle ~135°.

### 2.3 Other infrared/(sub)millimetre data

To refine the sampling of the spectral energy distribution (SED), we include photometric data, other than those obtained by Herschel and Spitzer, ranging from the IR to mm wavelength domain. At NIR and MIR wavelengths, we complement our data set of observations with flux measurements from the Wide-field Infrared Survey Explorer (WISE; Wright et al. 2010). Images belonging to the WISE all-sky survey at 3.4, 4.6, 12.1 and 22.2 µm were retrieved from the NASA/IPAC Infrared Science Archive. To measure fluxes from those images, all stars and background sources were masked in the field of the galaxy. Since the galaxy is located at the edge of the map, a determination of the background level from an annulus around the galaxy was impossible. A background value was determined from averaging the flux over 100 apertures taken randomly within a 30-arcmin field in this masked image. The size of the apertures...
was adapted to the FWHM of the PSF in every waveband with a diameter of $4 \times$ FWHM. The FWHM in the different WISE bands varies from 6.1, 6.4 and 6.5 arcsec at 3.4, 4.6 and 12.1 µm up to 12 arcsec in the longest 22.2-µm waveband (Wright et al. 2010). The mean background value was subtracted from the masked image. From this masked and background-subtracted image, we determined fluxes in all wavebands from aperture photometry. With the stellar bulge being a prominent emission feature at NIR wave-lengths and the MIR images having a dominant contribution from PAHs and very small grains in the star-forming disc of NGC 4565, the size of the aperture was adjusted to fit the size of the galaxy emitting in each waveband. Fluxes of the stellar emission from WISE at 3.4 and 4.6 µm are consistent with IRAC flux measurements at 3.6 and 4.5 µm (see Table 1). Also the WISE emission in the longer wavelength bands corresponds well with measurements by IRAS and MIPS at similar wavelengths (see Table 1).

IRAS flux measurements at 12, 25, 60 and 100 µm were adopted from the catalogue of IRAS observations of large optical galaxies (Rice et al. 1988). Also the KAO observed NGC 4565 (Engargiola & Harper 1992) in wavebands centred at 100, 160 and 200 µm (see Table 1). ISO observations at 170 µm (32.5 ± 4.9 Jy; Stickel et al. 2004) were not considered here due to aperture effects (see also Wiebe et al. 2009). With Akari measurements at 65, 90, 140 and 160 µm (1.17 ± 0.12, 3.99 ± 0.22, 11.85 ± 1.07 and 12.58 ± 2.39 Jy, respectively) being up to one order of magnitude lower in comparison to results from other airborne and space telescope facilities (IRAS, ISO, MIPS, KAO) in overlapping wavebands, we decided not to take into account Akari flux measurements for the analysis in this paper.

At mm wavelengths, IRAM observations at 1.2 mm constrain the emission from cold dust in those wavebands. Alton et al. (2004) reported a flux density of 55 mJy when integrating along the major axis profile of the 1.2 mm emission for a 20-arcsec beam. Assuming the 1.2 mm emission is spatially distributed similar to that at 250 µm, we estimate a total flux density of 99 mJy at 1.2 mm for NGC 4565 (see Table 1).

### 2.4 GALEX and optical data (0.15–0.89 µm)

Galaxy Evolution Explorer (GALEX) FUV/NUV and optical Sloan Digital Sky Survey (SDSS) ugriz data for NGC 4565 were recovered from the GALEX and SDSS archives and reduced following the instructions outlined in Cortese et al. (2012b). Final maps were reduced to obtain pixel sizes of 1.5 and 0.4 arcsec for GALEX and SDSS observations, respectively. GALEX FUV and NUV channels are characterized by a FWHM of 4.2 and 5.3 arcsec, respectively (Morrisey et al. 2007), whereas the shape of the PSF in SDSS maps depends on the seeing during sky exposures. GALEX images for NGC 4565 were obtained in a single run of 1693 s exposure time, with SDSS images covered in 54 s of observing time. The first two panels on the top row of Fig. 1 show the GALEX FUV and NUV images, respectively. Each GALEX image clearly shows the low level of star formation activity in the innermost regions in NGC 4565. This lack of direct ultraviolet (UV) emission from young stars in the inner 80-arcsec region in NGC 4565 was already identified by Engargiola & Harper (1992) and Ohta & Kodaira (1995), relating this deficiency in star formation to a bar structure in the central regions.

Fig. 1 presents the SDSS r-, g-, r-, i- and z-band images in the last two top panels and first three panels on the second row. The main emission features in the optical SDSS images are the stellar disc and bulge. At shorter optical wavelengths (in particular u band), a large fraction of the stellar light is extinguished along the line-of-sight by the prominent dust lane. Towards longer wavelengths, the dust obscuration diminishes and also the boxy shape of the bulge becomes more evident. This boxy shaped bulge is thought to result from a bar in the inner region of NGC 4565 viewed edge-on (Combes & Sanders 1981) and is therefore considered to be a pseudo-bulge. A genuine bulge component is also present in NGC 4565, but its spatial extent and intensity is inferior to the dominant emission from the pseudo-bulge (Kormendy & Barentine 2010).

Both from GALEX and SDSS images, the warping structures on the south-eastern and north-western sides of the galaxy’s disc can be discerned. Especially at UV wavelengths, the warps are prominent emission features, indicative for the ongoing star formation in these tidally disrupted parts of the ISM in NGC 4565. Flux densities for the stellar emission in UV/optical wavebands were adopted from Cortese et al. (2012b) for GALEX FUV/NUV and SDSS gri bands (see Table 1). Fluxes for the remaining SDSS bands ($u$, $z$) were obtained in a homogeneous way from aperture photometry similar to the procedure for the other SDSS bands.

### 3 DUST ENERGY BALANCE IN NGC 4565

#### 3.1 SKIRT AND FITSKIRT

SKIRT (Baes et al. 2003, 2011) is a 3D Monte Carlo radiative transfer code designed to model the absorption, scattering and thermal re-emission of dust in a variety of environments: circumstellar discs (Vidal & Baes 2007), clumpy tori around AGN (Stalevski et al. 2012) and different galaxy types (Baes et al. 2010; De Looze et al. 2010; Gadotti, Baes & Falony 2010; Gomez et al. 2010).

FITSKIRT (De Geyter et al. 2012) is a fitting routine that combines the output of SKIRT with a genetic algorithm optimization library to obtain a best-fitting model for the stellar and dust components in a galaxy. In this manner, the best-fitting parameters for a variety of models and stellar and dust geometries can be derived when providing an input image from which the stellar emission and dust obscuration effects can be modelled.

#### 3.2 Radiative transfer modelling

##### 3.2.1 Model I

A first model accounting for the observed optical properties of stars and dust in NGC 4565 is obtained from the FITSKIRT fitting algorithm. FITSKIRT was applied to determine the best-fitting model parameters using our SDSS g-band image (see Fig. 1, last top panel) as a reference image. We choose the g-band image for the fitting procedure, since the emission in this waveband is dominated by the old stellar population. At shorter wavelengths, the contribution from young stars increases rapidly. The g-band image is also optimal to perceive the attenuation effects of dust, which diminishes at longer wavelengths. Prior to the FITSKIRT fitting algorithm, the stellar warps on both ends of the galactic disc were masked in this image.

The input model for the stellar component consists of three components. The first component is an exponential disc with scale length $h_R$ and height $h_z$ to fit the disc component:

$$
\rho(R, z) = \rho_0 \exp \left( -\frac{R}{h_R} \right) \exp \left( -\frac{|z|}{h_z} \right).
$$

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The second component is a flattened Sérsic geometry for the outer bulge component. It is modelled as
\[
\rho(r) = \rho_0 S_n \left( \frac{m}{r_{\text{eff}}} \right),
\]
where \( m = \sqrt{R^2 + z^2} \) is the spheroidal radius (\( q \) is the intrinsic flattening or axial ratio) and \( S_n(s) \) represents the dimensionless 3D spatial density of a model that deprojects to a projected surface brightness profile:
\[
I(r_p) = I_0 \exp \left[ -b_n \left( \frac{r_p}{r_{\text{eff}}} \right)^{1/n} \right],
\]
where \( b_n \) is a dimensionless constant that can be approximated by the formula: \( b_n = 2n - \frac{1}{2} + \frac{3}{4}q^2 + \frac{1}{8}(1 + q^2)^{3/2} + 0.125q^3(1 + q^2)^{1/2} \), as derived by Ciotti & Bertin (1999). Finally, the third component consists of an inner bulge described as a Sérsic function (same as equation 2 but with \( m = \sqrt{R^2 + z^2} \)) with a Sérsic index \( n \sim 1.4 \) and a scale length \( r_{\text{eff}} \sim 1.3 \) arcsec, similar to the values reported in Kormendy & Barentine (2010). The luminosity of the inner bulge was scaled to a fixed value \( L_V \sim 3.6 \times 10^4 \) \( L_V/\odot \), optimized to reproduce the g-band minor axis profile (see Fig. 3, last bottom panel).

For the SED of the stars, we assume a Maraston (1998, 2005) single stellar population (SSP) parametrization. The age and metallicity of the SSP in the disc and bulge are fixed to 8-Gyr old and a solar metallicity (\( Z = 0.02 \)), as obtained from diagnostic plots of several Lick indices (see figs 5 and 6 in Proctor, Sansom & Reid 2000). The age of the SSP is assumed to be the same in the stellar disc and bulge, since the boxy bulge is considered to be a pseudo-bulge. The intensity of the stellar components is set by its V-band luminosity (\( L_V \)).

Earlier studies of the dust geometry in NGC 4565 have suggested a ring-like structure for a large fraction of the dust (Engargiola & Harper 1992; Ohta & Kodaira 1995; Neininger et al. 1996; Kormendy & Barentine 2010; Laine et al. 2010). Therefore, our dust model is composed of an exponential disc with a density distribution as represented in equation (1) and a ring with a Gaussian distribution in radial direction and an exponential vertical profile:
\[
\rho(R, z) = \rho_0 \exp \left[ -\frac{(R - R_0)^2}{2\sigma^2} \right] \exp \left[ -\frac{|z|}{h_z} \right],
\]
where \( R_0, \sigma \) and \( h_z \) describe the radius of the ring, the radial dispersion and the vertical scale height. The content of the dust reservoir in NGC 4565 is scaled by the dust mass in each of the dust components. The dust population is assumed to be uniform across the entire galaxy, consisting of a composition of dust particles with a fixed grain size distribution. The abundances, extinction and emissivity of the dust mixture are taken from the Draine & Li (2007) model for the dust in our own Galaxy.

With the different model components for the stars and dust in NGC 4565 and furthermore assuming a variable inclination for the galaxy, the degrees of freedom in the fitting procedure amount to 15 parameters. Our model does not account for the continuum emission from the AGN in the centre of NGC 4565, since the low [Ne II] 15.55 \( \mu \)m to [Ne II] 12.81 \( \mu \)m ratio (~1; Laine et al. 2010) suggests that the AGN continua are weak in this galaxy (Deo et al. 2007). Table 2 gives an overview of the best-fitting parameters obtained from the FITSKIRT algorithm. Since a sufficient number of photon packages were used to obtain this best-fitting model (a photon package is a collection of a number of photons, with the ensemble of all photon packages representing the radiation field in

| Geometry | \( i [\text{°}] \) | 87.25 |
|----------|----------------|-------|
| Stellar disc | \( h_R \) (arcsec) | 102 |
| | \( h_z \) (arcsec) | 8.2 |
| | \( L_V \) \( (L_V/\odot)^n \) | \( 4.9 \times 10^{10} \) |
| Outer bulge | \( n \) | 2.12 |
| | \( r_{\text{eff}} \) (arcsec) | 20.2 |
| | \( q \) | 0.744 |
| | \( L_V \) \( (L_V/\odot)^n \) | \( 2.0 \times 10^{10} \) |
| Dust disc | \( h_R \) (arcsec) | 223 |
| | \( h_z \) (arcsec) | 2.5 |
| | \( M_0 \) \( (M_V/\odot) \) | \( 6.3 \times 10^7 \) |
| Dust ring | \( R \) (arcsec) | 146 |
| | \( h_R \) (arcsec) | 4.3 |
| | \( \sigma \) (arcsec) | 41.2 |
| | \( M_0 \) \( (M_V/\odot) \) | \( 4.0 \times 10^7 \) |

\[ a \text{The luminosity of the stellar disc is later adjusted to } L_V = 4.0 \times 10^{10} \odot \text{ to correct for the emission from young stars in the V band when supplementing our model with star formation.} \]
the galaxy), we assume an uncertainty of about 20 per cent for all derived parameters (see also De Geyter et al. 2012).

Exploring the compatibility of this model with optical data, we compare the observed g-band image (see Fig. 3, top row) and the major- and minor axis g-band profiles for several galactocentric radii (see Fig. 3, bottom row) with the modelling results. The resemblance between modelled and observed images, major- and minor axis profiles suggests that our radiative transfer model is representative for the observed stellar emission and optical properties of dust in NGC 4565. For the inclination, we find a best-fitting value of $\sim 87.25$, which is in agreement with earlier reported results. Indeed, HI observations of the gaseous disc in the past have constrained the inclination $>84^\circ$ (Ruben 1991), with more recent HI observations reporting an inclination angle of 87.5 (Zschaechner et al. 2012). From a similar radiative transfer study, Alton et al. (2004) estimated an inclination of $88^\circ$.

From the stellar distribution in this radiative transfer model, we infer a bulge-to-disc ratio of B-to-D $\sim 0.47$ in the V band. The model indicates a total dust mass of $M_d \sim 1.0 \times 10^8 M_\odot$ of which about 60 per cent resides in an exponential disc, with the remaining 40 per cent distributed in a broad ring structure. The scale size of the exponential dust disc ($h_R \sim 223$ arcsec, $h_z \sim 2.5$ arcsec) agrees well with the results obtained from a similar radiative transfer modelling procedure in Alton et al. (2004) ($h_R \sim 200$ arcsec, $h_z \sim 4.3$ arcsec). The diffuse dust component in Alton et al. (2004) has a face-on optical depth $\tau_V \sim 0.63$, which corresponds to $M_d \sim 9.7 \times 10^7 M_\odot$, when assuming the same extinction cross-section for the Draine & Li (2007) dust composition applied in our dust model.

Figure 4 represents the SED as modelled by our radiative transfer code SKIRT (solid, black line). Available flux measurements were overlaid on this SED model to describe the observed SED. An overview of flux measurements for NGC 4565 is given in Table 1. Less accurate flux measurements were considered non-relevant for this analysis and were omitted from this plot (see Section 2). Although our model shows great resemblance with the g-band image and major- and minor axis profiles, the SED in Fig. 4 shows that our model fails to reproduce the emission from NGC 4565 in the UV and MIR/FIR wavelength domains. The higher observed UV and MIR emission from NGC 4565 indicates the presence of a young stellar population in this galaxy, which is not yet accounted for by our model.

### 3.2.2 Model 2: standard model with star formation

In an attempt to reproduce the UV and IR emission observed from NGC 4565, we add a star formation component to the model constructed in Section 3.2.1. Rather than using an empirical star formation template (e.g. MacLachlan et al. 2011), we adopt starburst templates from the library of pan-spectral SED models for the emission from young star clusters with ages $<10$ Myr presented in Groves et al. (2008). The SED templates in Groves et al. (2008) were generated from a one-dimensional dynamical evolution model of H II regions around massive clusters of young stars, the stellar spectral synthesis code STARBURST 99 (Leitherer et al. 1999) and the nebular modelling code MAPPINGS III (Groves 2004). The STARBURST99 templates used in these models correspond to an instantaneous burst with a Kroupa (2002) broken power law initial mass function (IMF). The main parameters controlling the shape of the emission spectrum from these young stellar clusters are the metallicity ($Z$) of the gas, mean cluster mass ($M_{cl}$), age and compactness ($C$) of the stellar clusters, the pressure of the surrounding ISM ($P_\text{ISM}$) and the cloud covering fraction ($f_{\text{PDR}}$). The metallicity is chosen similar to the metallicity of the old stellar population in NGC 4565 ($Z \sim 1 Z_\odot$). To eliminate the age parameter, we use age-averaged templates obtained from averaging spectra over 21 cluster ages, ranging from 0.01 to 10 Myr in steps of 0.5 Myr. For the mean cluster mass and ISM pressure, we assume fixed values $M_{cl} \sim 10^3 M_\odot$ and $P_\text{ISM} \sim 10^6$ cm$^{-3}$ K. Those approximations are justified since more massive star clusters can be simply thought of as the superposition of several individual clusters. A variation in the ISM pressure furthermore mainly affects the nebular emission lines rather than altering the shape of the emission spectrum (Groves et al. 2008). For the cloud covering fraction $f_{\text{PDR}}$, we will assume a value of $f_{\text{PDR}} = 1$ to describe the emission from heavily obscured star-forming regions in NGC 4565.
The choices for $M_\odot$ and $P/k$ leave us to explore the effect of the compactness parameter $C$ on the shape of the SED for values of $\log C$ ranging between 4.5 and 6.5. With the compactness parameter controlling the position of the FIR emission bump by shifting it in peak wavelength without significantly altering the width of the emission curve, the parameter $C$ is closely related to the dust temperature (Groves et al. 2008). Although the emission at 24 $\mu$m might still be dominated by hot dust emission predominantly heated by young stars of ages $<10$ Myr, the contribution of diffuse dust emission is non-negligible at FIR and submm wavelengths. Therefore, it is difficult to constrain the shape of the emission spectrum from hot dust grains and thus to determine the appropriate value of $\log C$ in the models. With values of 6.5 for $\log C$ describing extremely compact star-forming regions, which are characteristic for the high-pressure ISM and/or massive young clusters in starburst galaxies, we assume a more moderate compactness factor of $\log C = 5.5$. A value of $\log C = 5.5$ is furthermore consistent with the dust temperature $T_d = 30$ K found for the warm dust component heated predominantly by stellar emission originating from star-forming regions (Engargiola & Harper 1992).

In addition to these young stellar clusters, we have also included a component of ultracompact H II regions in our model. Those ultracompact H II regions are typically used to simulate the conditions during the earliest stages of the cluster lifetime ($<10^6$ yr) when the newly formed stars are likely still buried in their separate birth clouds and a typical star cluster consists of an ensemble of individual ultracompact H II regions, rather than providing a single source of radiation originating from the smoothed emission from all stars in the stellar cluster. The templates for these ultracompact H II regions were constructed in a similar way as for the dusty H II regions with their surrounding photodissociation regions (PDRs), but each birth cloud is modelled individually and the template for the entire compact H II regions is constructed by co-adding the individual SEDs (see also Groves et al. 2008 for a more thorough description of the models). The SED template for the ultracompact H II is jointed with the SED template for the H II regions with $f_{\text{FR}} = 1$, under the assumption that they produce new stars at the same rate. This joint SED template will be scaled until the observed emission at 24 $\mu$m is reproduced by our model.

Emission spectra for the young stellar population with ages between 10 and 100 Myr are taken from the STARBURST99 library (Leitherer et al. 1999). The STARBURST99 templates were constructed for an instantaneous burst with a solar metallicity, a burst mass $M_*=10^6 M_\odot$, and a Kroupa (2002) broken power-law IMF, which is consistent with the stellar emission spectra used to compute the SED of stellar clusters younger than 10 Myr in Groves et al. (2008). With the young stellar clusters $<10$ Myr emitting predominantly in MIR wavebands, the non-ionizing radiation of the evolved stellar population up to ages 50–100 Myr will dominate the UV emission spectrum in NGC 4565 (Calzetti et al. 2005).

Although most details on the exact location of star-forming complexes within the disc vanish along the line-of-sight, we try to recover the average geometry of unobscured and dust-enshrouded star-forming regions from the emission profiles along the major axis of NGC 4565 in the GALEX NUV and MIPS 24 $\mu$m wavebands (see Fig. 5, top and bottom panel, respectively). By allocating a certain distribution to the star-forming regions, the proper amount of attenuation due to the foreground screen of diffuse dust is applied on every specific location within the galaxy. Indeed, due to the high inclination angle the different star-forming regions in NGC 4565 will experience wide ranges in obscuration depending on whether they are located on the near- or far end of the galaxy’s disc with respect to the line-of-sight. From the major axis NUV emission profile, we derive a ring-like geometry for the unobscured star formation component, modelled as a ring (see equation 4) with central radius $R \sim 290$ arcsec and width $\sigma \sim 72$ arcsec. The scale height of the ring of young stars is assumed to be identical to the scale height of the stellar disc populated with the old stellar population ($h_* \sim 8.2$ arcsec). From the MIPS 24-$\mu$m image, it becomes immediately evident that a component of obscured star formation resides in a narrow ring coinciding with the CO molecular ring centred at 80–100 arcsec (Sofue & Nakai 1994; Neininger et al. 1996). The parameters for the ring (see equation 4) of embedded star formation are chosen similar to the molecular CO ring with central radius at $R \sim 90$ arcsec, width $\sigma \sim 3.4$ arcsec and scale height $h_* \sim 1.5$ arcsec. Aside from this concentration of embedded localized sources distributed in a ring, a more extended component of obscured star formation seems also present in NGC 4565. To reproduce the more extended 24 $\mu$m emission in the major axis profile, obscured star-forming complexes are distributed in the ring harbouring the unobscured star formation component and the dust ring containing $\sim 40$ per cent of the diffuse dust component in NGC 4565. From the 24 $\mu$m major axis profile (see Fig. 5, bottom panel), the presence of a central dust disc becomes evident. This central dust concentration is modelled as an exponential disc with radial and vertical scale lengths of 5 and 1 arcsec, respectively. The distribution of this central dust reservoir ($M_\text{dust} = 3 \times 10^4 M_\odot$) is truncated at a radius of 20 arcsec.
When fitting the UV and 24 μm emission of NGC 4565, we require an average star formation rate in the outer regions of the disc (with most of the young stars between galactocentric radii of 200 and 400 arcsec) of 3.1 M⊙ yr⁻¹ for the 10–100 Myr old stellar population. The heavily obscured star-forming regions of ages < 10 Myr seem to reproduce the observed 24 μm emission when assuming an average star formation rate of 2.2 M⊙ yr⁻¹ over the past 10 Myr. The dust-enshrouded star-forming clumps in the narrow ring produce young stars at a rate of SFR ~ 0.6 M⊙ yr⁻¹, with a somewhat higher productivity for the dense clouds distributed in the ring of unobscured star formation (SFR ~ 0.8 M⊙ yr⁻¹) and the diffuse dust ring (SFR ~ 0.8 M⊙ yr⁻¹). This obscured star formation rate inferred from the models is perfectly consistent with the SFR ~ 2.0 M⊙ yr⁻¹ calculated in Wu & Cao (2006) from the FIR luminosity L_FIR in NGC 4565, corrected for the contribution from an old stellar population (Satyapal et al. 2005). Also monochromatic SFR tracers such as the 24 μm emission result in similar estimates for the SFR ~ 2.0 M⊙ yr⁻¹ (Rieke et al. 2009).

Accounting for all young stellar objects in our model, we derive an average star formation rate of 5.3 M⊙ yr⁻¹ over the past 100 Myr. The small decrease in the star formation activity during the recent 10 Myr is in agreement with a delayed exponential law describing the star formation history in other late-type spiral galaxies (Boselli et al. 2001; Gavazzi et al. 2002). The total star formation activity in our model (SFR ~ 5.3 M⊙ yr⁻¹) is furthermore consistent with the total star formation rates derived from a combination of unobscured and obscured star formation tracers (e.g. SFR_{UV+24 μm} ~ 3.2 M⊙ yr⁻¹, Zhu et al. 2008; SFR_{UV+TIR} ~ 3.9 M⊙ yr⁻¹, Buat et al. 2005).

This model with both unobscured and obscured star formation is able to account for the young stellar emission at UV/MIR wavelengths (see images on the first and third rows of Fig. 6, respectively) and remains in agreement with the optical constraints (see images on the second row and plots on the last row of Fig. 6) upon reducing the luminosity of the old stellar population in the disc by a small factor to correct for the emission from those young stars in optical wavebands (L_V ~ 4.0 × 10^{10} L⊙). The total dust mass associated with the young stellar clusters (<10 Myr) in our models is M_d ~ 3.6 × 10^7 M⊙, which increases the total dust mass in NGC 4565 to M_d ~ 1.4 × 10^8 M⊙. The geometry for the embedded, localized sources in NGC 4565 is able to account for the 24 μm emission in NGC 4565 as observed from the major axis profile (see Fig. 5, bottom panel) or the MIPS 24-μm image (see images on the fourth row of Fig. 6). The geometry for the less obscured young stars with ages >10 Myr is capable of reproducing the emission profile along the major axis (see Fig. 5, top panel) as well as the emission originating from the entire galaxy at NUV wavelengths (see Fig. 6, third row). The peak at R ~ 420 arcsec in the NUV major axis profile belongs to the residual star formation in the warp structure on the north-western side of the galaxy’s disc and is not taken into account in our model. The stellar clusters do have a more clumpy distribution throughout the disc than was accounted for by our model.

More critical is the inconsistency in the FIR/submm, where our radiative transfer model continues to underestimate the observed emission spectrum for NGC 4565 (see Fig. 7). Even though the model accounts for the dust heating provided by star formation, the modelled dust emission underestimates the observed FIR and...
submm emission by a factor of $\sim 3$–4, in correspondence to dust energy balance studies for other nearby edge-on spirals.

At mid-IR wavelengths, our model also underestimates the observed emission originating mainly from PAHs in those wavebands by a factor of $\sim 2$ (see Fig. 7). This dissimilarity might on the one hand be a consequence of the fixed PAH template used in the models to construct the emission spectra reported in Groves et al. (2008). Based on the average variation by a factor of 2 obtained from a study of the mid-IR emission for a large sample of nearby galaxies (Smith et al. 2007), we should interpret this off-set as a diagnostic for the specific ISM conditions in NGC 4565 rather than attributing it to a deficiency of star formation in our model. With the PAH template in Groves et al. (2008) being optimized to fit the Spitzer IRS observations of two interacting galaxy pairs (NGC 4676 and NGC 7252), small deviations in ISM conditions and interstellar radiation field hardness might explain the difference in observed PAH emission for NGC 4565. On the other hand, the diffuse emission from cold dust at FIR/submm wavelengths has not yet properly been accounted for by our radiative transfer model and might be responsible for an important fraction of the PAH emission in galaxies (e.g. Haas, Klaas & Bianchi 2002; Bendo et al. 2008).

4 DISCUSSION

4.1 Dust energy balance

Although in agreement with the observed properties of stars and dust at optical wavelengths and the young stellar emission in UV and MIR wavebands, the model from Section 3.2.2 underestimates the observed dust emission in the FIR and (sub)millimetre wavelength regime. Whereas the dust emissivity of dust grains could influence the shape of the SED at some wavelengths (Alton et al. 2004), the inconsistency in the integrated dust energy balance cannot be solely caused by a grain population with dust emissivities diverging from Galactic values at longer wavelengths. Indeed, the multifrequency sampling of the dust SED allows us derive a true discrepancy in the integrated dust energy balance, which is impossible to solve by assuming divergent FIR emissivities due to a different grain composition and structure compared to the typical Galactic dust reservoir. Although we cannot exclude a possible variation in dust emissivity at submm/mm wavelengths (diverging from the assumed emissivity in the Draine & Li 2007 dust model), we argue that this discrepancy is likely due to a clumpy morphology of the dust in NGC 4565 unresolved in the currently available data set of FIR/submm observations for NGC 4565. Because of the high density in those clumps, they hardly contribute to the attenuation of stellar light and can therefore not be distinguished from optical observations.

Because of the edge-on inclination of NGC 4565 and the current limitations in resolving power of IR instrumentation, we are not capable of constraining the exact location and size of those dust clumps and rather have to remain content with information on the average mass and temperature of this clumpy dust reservoir. Those dust clumps either correspond to quiescent clouds where star formation has not (yet) been initialized. On the other hand, this extra dust reservoir might reside in the outer shells of star-forming complexes where the heating by the young embedded objects becomes negligible due to the high optical depth of the dust shells in the immediate vicinity of the stellar source.

To estimate the amount of dust residing in clumps, we fit the DustEm dust model with a Draine & Li (2007) grain composition to the residual fluxes after subtracting the dust emission in our model from the observed flux densities (IRAC 5.8, 8.0 $\mu$m, WISE 12.1 $\mu$m, IRAS 12 and 100 $\mu$m, Kuiper 100 and 160 $\mu$m, MIPS 160 $\mu$m, SPIRE 250, 350, 500 $\mu$m and IRAM 1.2 mm). Fig. 7 shows the resulting SED (black, solid curve) when supplementing our

Figure 7. Model 2: the modelled SED as obtained with the SKIRT model (dashed purple line) overlaid with the observed fluxes (see Table 1). Model 2 corresponds to model 1 supplemented with young star clusters ($<100$ Myr). To estimate the dust mass residing in quiescent dust clumps in NGC 4565, we fit the DustEm dust model with a Draine & Li (2007) grain composition (green, dotted line) to the residual fluxes for NGC 4565 in the FIR and submm wavebands. The model with star formation and an additional dust reservoir distributed in clumps can explain the multiwavelength observations for NGC 4565 (see black, solid curve).
model with a dust reservoir of \( \sim 1.5 \times 10^8 \text{ M}_\odot \) at a temperature of \( T_d \sim 16.9 \text{ K} \) (green, dotted curve). The average temperature \( T_d \sim 16.9 \text{ K} \) obtained for the residual dust reservoir is consistent with the temperature of the diffuse dust in the disc of NGC 4565 (\( T_d \sim 14–18 \text{ K} \)) and the average colour temperatures derived in other nearby spiral galaxies at wavelengths longwards of 160 \( \mu \text{m} \) (e.g. M81, M83 and NGC 2403; Bendo et al. 2012a). At mid-IR wavelengths, this additional dust reservoir is capable of explaining the PAH emission in NGC 4565. Also at FIR/submm wavelengths, a massive cold dust component seems adequate to explain the emission in those wavebands. At 60 and 70 \( \mu \text{m} \), our radiative transfer model however overestimates the observed fluxes. This inconsistency could be an indication for the star-forming complexes in NGC 4565 to be more compact than currently accounted for by our model. Because of the high inclination angle of NGC 4565, a direct identification of those complexes is heavily obscured star-forming complexes within the disc of the galaxy and their individual size distribution and compactness is impeded.

With these additional cold dust clumps, the total dust mass in NGC 4565 amounts to \( M_d \sim 2.9 \times 10^8 \text{ M}_\odot \) of which one-third has a diffuse distribution throughout the disc and the remaining two-thirds resides in clumps. From those dense dust clouds, the majority (\( \sim 80 \) per cent) is not heated by UV-emitting sources while the remainder (\( \sim 20 \) per cent) host embedded star-forming complexes. The total dust mass \( M_d \sim 2.9 \times 10^8 \text{ M}_\odot \) in our model is consistent with the dust reservoir accounted for by BLAST observations (Wiebe et al. 2009), but \( \sim 10 \) times higher than the dust content reported in Engargiola & Harper (1992) (\( 2.7 \times 10^7 \text{ M}_\odot \)) based on IRAS observations up to 100 \( \mu \text{m} \). This confirms the need for longer wavelength observations to trace the entire dust reservoir in a galaxy (e.g. Bendo et al. 2010, 2012a; Gordon et al. 2010).

The solution to the dust energy budget problem for NGC 4565 accounting for a sizable fraction of the dust (\( \sim 65 \) per cent) distributed in compact clumps is in agreement with other energy balance studies of edge-on objects (e.g. Bianchi 2008; Baes et al. 2010; Popescu et al. 2011; De Looze et al. 2012; Holwerda et al. 2012, etc.).

4.2 Gas-to-dust ratio

From the \( \text{H}I \) mass \( M_{\text{HI}} \sim 1.8 \times 10^{10} \text{ M}_\odot \) (Zschaechner et al. 2012) and the CO observations (\( M_{\text{HI}} \sim 2.9 \times 10^9 \text{ M}_\odot \)) presented in Neininger et al. (1996) (both quantities are already converted to our adopted distance of \( D = 16.9 \text{ Mpc} \)), we derive a total gas mass of \( \sim 2.9 \times 10^{10} \text{ M}_\odot \) for NGC 4565 when applying a correction factor of \( \sim 1.4 \) to correct for heavier elements. Based on the total dust mass in our radiative transfer model (\( M_d \sim 2.9 \times 10^8 \text{ M}_\odot \)), we estimate an average gas-to-dust ratio of \( \sim 100 \) in NGC 4565, which is in fair agreement with the gas-to-dust fraction found in other late-type spirals (e.g. \( \sim 120 \) in our own Galaxy; Zubko, Dwek & Arendt 2004) and the dust-to-\( \text{H}I \) mass fraction derived for Sb galaxies with normal \( \text{H}I \) components (as opposed to \( \text{H}I \) deficient objects) in the HRS sample of galaxies (Cortese et al. 2012a).

4.3 Dust heating sources

Investigating the dust heating in NGC 4565, we find that about 70 per cent of the FIR luminosity is attributed to the absorption of stellar light from the old stellar population, with the heating for the remaining 30 per cent provided by the young stellar population. In this calculation, we assume that the high-density dust clumps distributed throughout the diffuse ISM in NGC 4565 are merely heated by the old stellar population, since the temperature of the dust (\( T_d \sim 16.9 \text{ K} \)) in those dense dust cores does not suggest any link to localized embedded sources and is more closely related to the average temperature of the diffuse dust in the disc of the galaxy (\( T_d \sim 14–18 \text{ K} \)). Those heating fractions are in agreement with similar studies of dust heating mechanisms in other nearby galaxies, where the old stellar population also accounts for an important fraction of the FIR radiation (e.g. Bianchi 2008; Bendo et al. 2010, 2012a; Bouquen et al. 2011; Popescu et al. 2011). Similar results were obtained from accurate modelling of the absorption of stellar light as well as careful investigations of UV observations (Buat & Xu 1996; Charlot & Fall 2000; Witt & Gordon 2000; Cortese et al. 2008). Engargiola & Harper (1992) report a dust heating efficiency of 54 and 46 per cent for the young and old stellar population in NGC 4565, respectively. Considering that their IRAS observations only covered a wavelength range from 12 to 100 \( \mu \text{m} \), it is not surprising that they underestimated the contribution from the old stellar population to the dust heating with the more evolved stars contributing the bulk of dust heating at longer wavelengths.

Although the fraction of dust heated by the young stellar population in NGC 4565 is overall small (\( \sim 30 \) per cent), the contribution of hot dust emission from star-forming complexes is about 90 per cent in the 24-\( \mu \text{m} \) waveband, after which it decreases quickly to about 50 per cent at 70 \( \mu \text{m} \), 12 per cent at 160 \( \mu \text{m} \), less than 10 per cent in all SPIRE wavebands and below 5 per cent at 1.2 mm. Those values are consistent with the decreasing contribution of star-forming regions for increasing wavelength in M81, M83 and NGC 2403 (Bendo et al. 2012a), with a relative contribution of the young stellar population to the total dust heating of more than 50 per cent shortwards of 250 \( \mu \text{m} \) and dropping below 30 per cent in the SPIRE wavebands.

This self-consistent analysis of the dust heating mechanisms in NGC 4565 based on a realistic radiative transfer model confirms earlier studies reporting on a non-negligible fraction of the dust heating in spiral galaxies powered by the more evolved stellar population. With the young stellar population being responsible for less than half of the total IR dust emission originating from NGC 4565, we believe caution is needed when using the total IR emission in galaxies to trace the star-forming activity. Although a tight correlation was found between the star formation and the total IR dust emission for large samples of nearby spiral galaxies (Devereux & Young 1990; Devereux, Jacoby & Ciardullo 1995; Buat & Xu 1996; Kennicutt et al. 2009), this relation might be the result of an indirect link between the star-forming activity in galaxies and the total IR emission in galaxies, tracing the surface density of gas in the ISM. In the latter case, the SFR–\( L_{\text{IR}} \) correlation is governed by the Schmidt law (Schmidt 1959; Kennicutt 1998), relating the surface density of gas and star formation in galaxies, rather than the dust heating provided by star-forming regions.

5 CONCLUSIONS

We present a full radiative transfer analysis of the edge-on spiral galaxy NGC 4565, accounting for the absorption and scattering of stellar light by dust grains and its thermal re-emission at IR wavelengths. From a radiative transfer model fitting procedure to the optical SDSS g-band image, we determine the best-fitting parameters for the old stellar population and diffuse dust component in NGC 4565. To account for the observed UV and MIR emission in NGC 4565, we supplement our model with a young stellar population of age <10 Myr (\( \text{SFR} \sim 2.2 \text{ M}_\odot \text{ yr}^{-1} \)) and ranging from 10 to 100 Myr (\( \text{SFR} \sim 3.1 \text{ M}_\odot \text{ yr}^{-1} \)). The distribution of star-forming
complexes within the disc of the galaxy is constrained from major axis NUV and 24 µm emission profiles. Even though this young stellar population provides an additional power source for dust heating and hereby boost the emission at MIR wavelengths, the emission observed at wavelengths longwards of 100 µm remains underestimates by a factor of 3–4. This inconsistency in the dust energy budget of NGC 4565 suggests the presence of a sizable fraction (two-thirds) of the total dust reservoir (\( M_d \sim 2.9 \times 10^9 M_\odot \)) in a clumpy distribution with no associated young stellar sources. The distribution of those dense dust clouds would be such that they remain unresolved in current FIR/submm observations and hardly contribute to the attenuation by dust at optical wavelengths. The contribution from a grain composition with dust emissivity properties diverging from the typical Galactic standards can however not be discarded as a factor of influence on the high FIR/submm fluxes. The majority of dust heating in NGC 4565 is provided by the old stellar population (70 per cent), with the remaining dust heating powered by localized embedded sources in NGC 4565.

The results from this detailed dust energy balance study in NGC 4565 accords with similar analyses of other edge-on spirals, concluding that a significant fraction of the dust is distributed in such a way that the influence on the overall extinction of a galaxy is negligible but does emit strongly in FIR/submm wavebands.

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