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The thermal conductivity of the Earth’s core and implications for its thermal and compositional evolution

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Summary
Precise determinations of the thermal conductivity of iron alloys at high pressures and temperatures are essential for understanding the thermal history and dynamics of the metallic cores of the Earth. We review relevant high-pressure experiments using a diamond-anvil cell and discuss implications of high core conductivity for its thermal and compositional evolution.

Main text
The thermal conductivity of iron alloys is a key to understanding the mechanism of convection in the Earth’s liquid core and its thermal history. The Earth’s magnetic field is formed by a dynamo action that requires convection in the liquid core. Present-day outer core convection can be driven by the buoyancy of light element enriched liquid that is released upon inner core solidification in addition to thermal buoyancy associated with secular cooling. In contrast, before the birth of the inner core, the core heat loss must be more than the heat conducted down the isentropic gradient in order to drive convection by thermal buoyancy alone, which can be tight constraints upon the core thermal evolution.

Recent mineral physics studies throw traditional value of the Earth’s core thermal conductivity into doubt (Fig. 1). Conventionally the thermal conductivity of the outer
core had been considered to be $\sim 30 \text{ W m}^{-1}\text{ K}^{-1}$ that was estimated based on shock experiments and simple physical models including the Wiedemann-Franz law: $\kappa_{el} = LT\rho^{-1}$, where $\kappa_{el}$, $L$, $T$, and $\rho$ are electronic thermal conductivity, Lorenz number, temperature, and electrical resistivity, respectively [1]. Such relatively low core conductivity indicates that liquid core convection could have been driven thermally even with relatively slow cooling rate. However, in 2012–2013, our conventional view was challenged by both computational and experimental studies showing much higher core conductivity [2-4].

Since then, experimental determinations of the thermal conductivity of iron and alloys have been controversial (Fig. 1). Ohta et al. [5] measured the electrical resistivity of iron under core conditions in a laser-heated diamond-anvil cell (DAC). The results demonstrate relatively high thermal conductivity of $\sim 90 \text{ W m}^{-1}\text{ K}^{-1}$ for liquid Fe-Ni-Si alloy based on their measured resistivity for pure iron, Matthiesen’s rule, and Wiedemann-Franz law, which is compatible with \textit{ab initio} simulations [2,4]. On the other hand, flash laser-heating and fast thermal radiation detection experiments demonstrated the low core conductivity of 20–35 W m$^{-1}$K$^{-1}$ based on finite element method simulations [6,7], in accordance with traditional estimate [1]. Since transport properties that describe non-equilibrium phenomena are difficult to measure, the fact that determinations of the iron conductivity under core conditions became viable these days is remarkable success in mineral physics. Nevertheless, the discrepancy in core conductivity makes a big difference in the expected age of the inner core, mechanism of liquid core convection, and thermal history [3].

Despite a number of following studies based on a variety of different techniques, we still see a dichotomy of proposed core conductivity values (Fig. 1). The “saturation” resistivity, which is derived from the fact that the mean free path of electron-phonon interaction cannot be longer than the interatomic distance, gives the lower bound for conductivity. Such saturation resistivity lies between two clusters of reported high and low resistivity values. While the resistivity saturation is important in highly resistive transition metals and their alloys [3,8] (Fig. 2), the conventional estimate [1] did not include the effect of saturation in their models, which resulted in much higher resistivity than the saturation value and hence low core conductivity. The core electrical resistivity measured by recent DAC experiments [3,5,9] show resistivity saturation (Fig. 2), demonstrating the high core conductivity as far as the Wiedemann-Franz law holds with.
ideal Lorenz number (Fig. 1). Additionally, since temperature has a large effect on resistivity, temperature gradient in a laser-heated sample is an issue. An internally-resistance-heated DAC provides homogenous and stable sample heating and is thus a promising technique for conductivity measurements at high pressure and temperature \((P-T)\) [9]. The validity of the Wiedemann-Franz law under extreme conditions has been also an issue. Simultaneous measurements of the electrical resistivity and the thermal conductivity of iron alloy under core high \(P-T\) conditions will provide the decisive evidence for it.

As introduced above, the most recent high \(P-T\) measurements for Fe containing 2, 4, 6.5 wt.% Si using an internally-resistance-heated DAC have demonstrated that the thermal conductivity of Fe-12.7 wt.% (22.5 at.%) Si is \(\sim 88 \text{ W m}^{-1}\text{K}^{-1}\) at CMB conditions when the effects of resistivity saturation, melting, and crystallographic anisotropy at measurements are taken into account [9] (Fig. 1). Thermal conductivity of Fe-10 at.% Ni-22.5 at.% Si alloy, a possible outer core composition, could be \(\sim 79 \text{ W m}^{-1}\text{K}^{-1}\) considering the impurity effect of Ni [10]. Si exhibits the largest “impurity resistivity”, indicating that the \(79 \text{ W m}^{-1}\text{K}^{-1}\) is the lower bound for the thermal conductivity of the Earth’s liquid core. The core thermal evolution models by Labrosse [11] demonstrated that if liquid core convection has been driven by thermal buoyancy with the core thermal conductivity of \(79 \text{ W m}^{-1}\text{K}^{-1}\) at the CMB and no radiogenic heating in the core, the CMB temperature is calculated to be \(\sim 5500 \text{ K}\) at 3.2 Ga and \(\sim 4800 \text{ K}\) at 2.0 Ga. Such high CMB temperature suggests that the whole mantle was fully molten until 2.0–3.2 Ga. It is not consistent with geological records, calling for a different mechanism of core convection.

Chemical buoyancy may be an alternate means of driving convection in the core from the early history of the Earth. It has been proposed that the compositional buoyancy in the core could arise from the exsolution of MgO, SiO\(_2\), or both [12-14]. Recent core formation models based on the core-mantle distributions of siderophile elements suggested that core metals segregated from silicate at high temperatures, typically at \(3000–4000 \text{ K}\) and possibly higher [13,15], which enhances the incorporation of lithophile elements including Si and O, and possibly Mg into metals. It is suggested that the (Si,O)-rich liquid core may have become saturated with SiO\(_2\) upon secular cooling [14]. Indeed, the original core compositions proposed in recent core formation models include Si and O beyond the saturation limit at CMB conditions [15],
i.e. 136 GPa and 4000 K, leading to SiO$_2$ crystallization [13]. The rate of SiO$_2$ crystallization required to sustain geodynamo is as low as 1 wt.% per 10$^9$ years, which corresponds to a cooling rate of 100–200 K Gyr$^{-1}$ [14]. The most recent model of the core compositional evolution by Helffrich et al. [13] showed that MgO saturation follows SiO$_2$ saturation only when >1.7 wt.% Mg in the core. If this is the case, in addition to solid SiO$_2$, (Mg,Fe)-silicate melts exsolve from the core and transfer core-hosted elements such as Mo, W, and Pt to the mantle. The core-derived silicate melts may have evolved toward FeO-rich compositions and now represent the ultra-low velocity zones above the CMB.

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Figure 1. (a) Electrical resistivity and (b) thermal conductivity values at the top of the Earth’s core in the literature [1,2,4-7,9,16]. Filled symbols were calculated on the basis of the Wiedemann-Franz law with ideal Lorenz number ($L_0 = 2.44 \times 10^{-8}$ W Ω K$^{-2}$). Gray bands indicate (a) the range of saturation resistivity [9] and (b) thermal conductivity computed from the saturation resistivity and the Wiedemann-Franz law.
Figure 2. Temperature response of the electrical resistivity of (a) fcc iron estimated at 1 bar [8] (blue curve) and (b) hcp iron at 115 GPa [5]. Red curve and black line with gray uncertainty band indicates the predicted resistivity based on the Bloch-Grüneisen model with and without the resistivity saturation, respectively.