HIGH-RESOLUTION STUDY OF THE CLUSTER COMPLEXES IN A LENSED SPIRAL AT REDSHIFT 1.5: CONSTRAINTS ON THE BULGE FORMATION AND DISK EVOLUTION

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ABSTRACT
We analyze the clump population of the spiral galaxy Sp 1149 at redshift 1.5. Located behind the galaxy cluster MACS J1149.5+2223, Sp 1149 has been significantly magnified allowing us to study the galaxy on physical scales down to ∼100 pc. The galaxy cluste frame is among the targets of the Cluster Lensing And Supernova survey with Hubble (CLASH), an ongoing Hubble Space Telescope (HST) Multi-Cycle Treasury program. We have used the publicly available multi-band imaging data set to reconstruct the spectral energy distributions of the clumps in Sp 1149, and derive, by means of stellar evolutionary models, their physical properties. We found that 40% of the clumps observed in Sp 1149 are older than 30 Myr and can be as old as 300 Myr. These are also the more massive (luminous) clumps in the galaxy. Among the complexes in the local reference sample, the star-forming knots in luminous blue compact galaxies could be considered progenitor analogous of these long-lived clumps. The remaining 60% of clumps have colors comparable to local cluster complexes, suggesting a similar young age. We observe that the Sp 1149 clumps follow the $M \propto R^2$ relation similar to local cluster complexes, suggesting similar formation mechanisms although they may have different initial conditions (e.g., higher gas surface densities). We suggest that the galaxy is experiencing a slow decline in star formation rate and a likely transitional phase toward a more quiescent star formation mode. The older clumps have survived between 6 and 20 dynamical times and are all located at projected distances smaller than 4 kpc from the center. Their current location suggests migration toward the center and the possibility of being the building blocks of the bulge. On the other hand, the dynamical timescale of the younger clumps is significantly shorter, meaning that they are quite close to their birthplace. We show that the clumps of Sp 1149 may account for the expected metal-rich globular clusters population usually associated with the bulge and thick disk components of local spirals.

Key words: galaxies: evolution – galaxies: high-redshift – galaxies: individual (Sp 1149) – galaxies: star formation

Online-only material: color figures

1. INTRODUCTION
Recent results suggest that galaxies that are still star-forming at the present time (redshift, $z \sim 0$), e.g., late-type galaxies such as spirals have assembled more than 80% of their stellar mass between $0 \lesssim z \lesssim 2$ (e.g., Moster et al. 2013; Yang et al. 2012; Leitner 2012). It is well known that the cosmic star formation rate (SFR) density shows a peak at $z \sim 2$ and a clear decline between $0 \lesssim z \lesssim 2$ (see Lilly et al. 1996; Madau et al. 1996; Bouwens et al. 2011 for one of the updated diagrams of the cosmic SFR as a function of $z$). More recently it has become clear that SFR at all $z$ is probably dominated by normal, $L^*$ star-forming galaxies (Noeske et al. 2007; Wuyts et al. 2011). From those two results, it follows that normal spiral galaxies have strongly declining SFRs between $z = 2$ and the present.

At high $z$, galaxy morphologies do not follow the classical Hubble classification. Galaxies are intrinsically different and show features like very massive star-forming clumps that could hardly be associated with any star-forming region we observe in local galaxies. Using HST ultra-deep field data, Elmegreen et al. (2005) defined six morphological types, of which four (chain galaxies, clump clusters, double, tadpole) are based on the prominent presence of compact clumps, and the absence of any bulge or exponential profile. In particular, chain and clump cluster galaxies are probably the same class of objects but observed at different inclinations. At $z \gtrsim 1$, clumps are also observed in galaxies that show structural and kinematic properties typical of disks (e.g., Elmegreen et al. 2007; Jones et al. 2010; Genzel et al. 2011). It seems, therefore, that clump galaxies, in the absence of major mergers, will become the high-$z$ progenitors of modern disk galaxies (e.g., Elmegreen et al. 2009b). High-$z$ clumpy disks are observed to have high gas fractions, with spatially extended distributions that overlap the UV morphologies, and clear hints of rotation (e.g., Daddi et al. 2010; Tacconi et al. 2010). Observationally, there are hints that the Toomre (1964) $Q$ parameter of clumpy disks is $\lesssim 1$, suggesting that clumps form because of gravitational instability (e.g., Jones et al. 2010; Genzel et al. 2011). At the resolution achieved by these studies, the masses of the clumps are quite high ($10^7$–$10^9 M_\odot$), with radii between 0.5 and a few kiloparsecs. Theoretical models and simulations predict that cold accretion of gas could at the same time feed the galaxy prolonging the gas depletion time, assuming that the SFR stays constant, and induced instabilities in the gas to form clumps (see Elmegreen 2011a for a short review). In this scenario, the average surface density of SFR becomes a direct measure of the gas surface density present in the galaxy (Livermore et al. 2012). Under the assumption of a rotating and marginally unstable ($Q \sim 1$) disk, Livermore et al. (2012) argued that the mass of the more massive clumps that can form increases for higher disk
surface density. This finding can, indeed, explain why, when the properties of high-z clumps are compared to local star-forming regions, they appear to follow the same physical relations as the latter but at much higher scales of SFR, luminosity, total stellar mass, etc. (e.g., Swinbank et al. 2010; Tacconi et al. 2010).

A comparison study of the clump properties in clump clusters, chains, and young spirals shows that the former may evolve into spiral galaxies (Elmegreen et al. 2009a). Indeed, theoretical works and simulations show that clumps form in gravitationally unstable, highly turbulent disks (e.g., Bournaud et al. 2008; Ceverino et al. 2010). These clumps are expected to migrate toward the center of the system because of dynamical friction and tidal torques in about 10 times the dynamical time, i.e., a total time range of 0.5–1 Gyr. In the center clumps coalesce and slowly build up the bulge and the thick disk component (Bournaud et al. 2007; Elmegreen et al. 2008). However, this scenario works if clumps survive long enough to reach the center of the galaxy. In simulations and analytical studies, different initial conditions and assumptions lead to opposite fates of the clumps, i.e., they can be rapidly disrupted by stellar feedback (Murray et al. 2010) or survive to the initial gas loss for several dynamical times, so they could spiral inward (Krumholz & Dekel 2010). Streams of material coming from dissolving clumps could still migrate toward the center, but the bulge formation mechanism would be more inefficient (Genzel et al. 2011).

In spite of the progress made in understanding how galaxies have evolved from the time at which the cosmic SFR peaked (z ∼ 2) to the current decline, it is still challenging to get a complete census of the whole process. The major obstacles are the limited resolution and sensitivity, which create a bias toward galaxies with extreme SFRs and physical conditions not observed in the local universe. Recently, in order to study fainter galaxies with lower SFRs, researchers have focused their attention on galaxies that have been magnified by a strong gravitational lens, i.e., galaxy cluster potential well. This technique has enabled the study of galaxies and clump properties at spatial scales of a few hundreds parsecs, which are otherwise prohibitive (e.g., Swinbank et al. 2007; Stark et al. 2008; Jones et al. 2010; Livermore et al. 2012). The most important outcome of this high-resolution analysis is that in spite of the extreme conditions under which clumps and star formation happens in these galaxies, the underlying physical process that governs star formation is the same as observed in the local universe (e.g., Swinbank et al. 2010).

In this work, we will present an analysis of the clump population of a spiral galaxy at z ∼ 1.5. Following the nomenclature by Yuan et al. (2011), we will refer to it as Sp 1149. The system is located behind the galaxy cluster MACS J1149.5+2223 (hereafter MACS J1149), and thanks to gravitational magnification, allows us to resolve spatial scales down to ∼100 pc. According to Smith et al. (2009), Sp 1149 has a luminosity comparable to a local L* galaxy (Norberg et al. 2002), and is about 1 mag fainter than previously studied unmarshaled spirals at similar redshifts. We will try to constrain how star formation is proceeding in the galaxy by means of clump analysis. The current SFR of Sp 1149 is not extreme if compared to those of local spirals (in the range 1–6 M⊙ yr−1; Smith et al. 2009; Livermore et al. 2012). Yuan et al. (2011) reported clear signs of a rotationally supported disk and an increasing radial metallicity gradient toward the center. Therefore, Sp 1149 is a potential target for investigating the transition phase that extremely star-forming galaxies have experienced when evolving into modern late-type systems.

The paper is organized as follows. In Section 2, we introduce the data set and details of the photometric analysis. These data will be used to derive clump properties, e.g., age, mass, extinctions, effective radii, in Section 3. In Section 4, we introduce some new data, together with those from the literature, which will be used as a comparative sample, both at high z and in the local universe. Section 5 presents the analysis of the clump properties, in the context of local cluster complexes but also high-z clumps. Finally, in Section 6 we discuss the formation and fate of Sp 1149 clumps, evidence for the bulge formation, and the potential rate of metal-rich globular cluster progenitors. A discussion and summary will follow in the final section.

Throughout the paper we assume the following cosmological parameters: H0 = 73 km s−1 Mpc−1, ΩM = 0.27, and ΩΛ = 0.73.

2. DATA SET AND PHOTOMETRIC CATALOG

Our target is a lensed spiral galaxy with spectroscopically confirmed redshift z = 1.49 (Smith et al. 2009). Four multiple images of Sp 1149 have been identified in the galaxy cluster frame of MACS J1149. In this work we will focus on the tangential image A1.3 from Smith et al. (2009) shown here in Figure 1. The particular geometry of the gravitational potential causes this image of Sp 1149 to be almost equally lensed in both linear directions, thus avoiding the distortion effects that plague many other lensed galaxies. Using the lens model derived by Smith et al. (2009), we measured a linear magnification in the center of the galaxy of 2.7 × 3.1 (almost symmetric) at a position angle of 37° (Livermore et al. 2012; R. C. Livermore et al., in preparation). Therefore, we assume a uniform magnification of μ = 8 across the galaxy because the system is distant enough from the z = 1.5 critical line to not be affected by strong magnification gradients. The uncertainties on the magnification as a function of the position in the galaxy are smaller than the uncertainties associated with the computed galaxy cluster magnification map, i.e., ∼18%. Because the fluxes of each clump are scaled by the same magnification factor, this uncertainty will not change the shape of the clump spectral energy distribution (SED) but only the derived mass.

The galaxy cluster MACS J1149 has been included among the 25 targets of the Cluster Lensing And Supernova survey with Hubble (CLASH), an ongoing Hubble Space Telescope (HST) Multi-Cycle Treasury program (12065; PI: Marc Postman). MACS J1149 has been imaged in 17 HST filters from the UV to the NIR. In Table 1, we summarize the filters used in this analysis (11 out of 17), the aperture correction applied, and the photometric catalog of the extracted clump complexes in the ABmag system.

The CLASH team has released two reduced data sets of MACS J1149 with different final pixel scale resolutions, 0′03 pixel−1 and 0′065 pixel−1. The images have all been rescaled to the same pixel scale, aligned, and registered (see Postman et al. 2012 for more information). To carry out our analysis we use the set with 0′065 pixel−1 scale resolution. The latter is a good compromise between the finer resolution of the ACS/WFC and WFC3/UVIS cameras (0.05 and 0′04 pixel−1, respectively) and the worse scale sampling of the WFC3/IR (0′13 pixel−1).

To perform the photometric analysis we selected only the filters where the signal-to-noise ratio (S/N) was good enough

5 Released data available at http://archive.stsci.edu/prepds/clash/.
to detect the clumps (e.g., in the F225W, F275W, and F336W filters the galaxy is barely detected, so they are excluded from the analysis) and the filter width was such that it carried new information (e.g., the F625W has been excluded because it overlaps with the F555W, F606W, and F775W transmission). Of each large frame, we cut out a 13 arcsec$^2$ image centered on the spiral. The Advanced Camera for Surveys (ACS) frames have then been combined with the IRAF task imcombine, using the function SUM to improve the S/N of the sources. This combined frame has only been used to identify the clumps by eye. To do aperture photometry, we used a radius of 3 pixel ($\sim 0.2$) and a sky annulus 1 pixel wide located at a distance of 4 pixel. The aperture radius has been chosen large enough to contain the stellar FWHM of the IR camera. The combined frame used for clump detection and the aperture size applied to do photometry is shown in Figure 9.

As we will discuss in Section 3.2, some of the clumps are extended, others unresolved or barely resolved. To compensate for the missing flux caused by the use of a finite aperture radius we used unsaturated stars detected in each initial MACS J1149 frame to estimate the aperture correction. In the case of extended sources this correction factor is a lower limit to the fraction of missed flux. In other words, we have only corrected for the missing flux assuming a stellar point-spread function (PSF). The photometry has also been corrected for Galactic extinction (Schlegel et al. 1998).

Table 1 contains the final photometric catalog of the 31 clumps. These fluxes have been scaled down by the factor $\mu = 8$ to take into account the magnification of the lens. Photometry with a magnitude error larger than 1 mag has been discarded and considered as non-detection (empty cells in the table).

3. SP 1149 CLUMP ANALYSIS

3.1. Determining Age and Mass of the Cluster Complexes

Multi-band photometry has been used to perform an SED analysis of each clump.

The evolutionary tracks have been produced with the publicly available Yggdrasil models$^6$ (Zackrisson et al. 2011). We assume two different metallicities $Z = 0.004$ and 0.008, Starburst99 single stellar population (SSP) models based on Padova asymptotic giant branch tracks (Leitherer et al. 1999; Vázquez & Leitherer 2005), and a universal Kroupa (2001) initial mass function (IMF) in the interval 0.1–100 $M_\odot$. To regulate the relative contribution from nebular emission to the final SED we assumed a covering factor, $f_{\text{cov}} = 1$. This assumption also determines the Lyman continuum escape fraction, $f_{\text{esc}}$ through the relation $f_{\text{esc}} = 1 - f_{\text{cov}}$, i.e., none in our case. This choice of parameters seems reasonable if one considers that the radius inside which we do photometry corresponds to a projected physical size of about 600 pc, i.e., large enough to have a high probability for the ionizing photons to be absorbed. Finally, we have simulated three star formation histories (SFHs), instantaneous burst (e.g., SSP), constant SFR for 10 Myr, and for 100 Myr (i.e., it produces a burst of star formation with constant SFR lasting a given interval of Myr, after which the SFR drops to zero).

The $\chi^2$-algorithm used is described in Adamo et al. (2010). In short, comparing models and observed fluxes, the algorithm gives a best solution for a combination of age, internal extinction, and mass. We allowed data points with photometric errors as large as 1 mag to be included in the fit (note that the weight on the fit of these data is, therefore, very small). We have implemented the algorithm with the analysis of 68% confidence limits on the determination of the clump parameters (associated errors in Table 2). According to the prescription given by Lampton et al. (1976), the uncertainties associated with the best outcome parameters can be determined if the number of fitted data points is larger than the number of parameters (e.g., three in our case). In our catalog, only source 21 did not pass this criterion (detected in only three passbands), therefore we could not perform the error analysis.

$^6$ http://ttt.astro.su.se/~ez
### Table 1

**Photometric Catalog**

| ID | F390W | F475W | F555W | F606W | F775W<sup>a</sup> | F814W<sup>b</sup> | F850LP<sup>c</sup> |
|----|-------|-------|-------|-------|-----------------|-----------------|-----------------|
| Camera | WFC3 | ACS | ACS | ACS | ACS | ACS | ACS |
| UVIS | WFC | WFC | WFC | WFC | WFC | WFC | WFC |
| Exp. time (s) | 4781 | 4136 | 9000 | 4128 | 4094 | 13548 | 8280 |
| AC<sup>d</sup> | −0.15 | −0.16 | −0.18 | −0.16 | −0.22 | −0.33 | −0.25 |

| ID | F105W<sup>e</sup> | F110W<sup>f</sup> | F125W<sup>g</sup> | F160W<sup>h</sup> |
|----|-----------------|-----------------|-----------------|-----------------|
| Camera | WFC3 | WFC3 | WFC3 | WFC3 |
| IR | IR | IR | IR |
| Exp. time (s) | 2814.7 | 2414.7 | 2514.7 | 5029.3 |
| AC | −0.38 | −0.39 | −0.45 | −0.60 |

| ID | F105W<sup>e</sup> | F110W<sup>f</sup> | F125W<sup>g</sup> | F160W<sup>h</sup> |
|----|-----------------|-----------------|-----------------|-----------------|
| Camera | WFC3 | WFC3 | WFC3 | WFC3 |
| IR | IR | IR | IR |
| Exp. time (s) | 2814.7 | 2414.7 | 2514.7 | 5029.3 |
| AC | −0.38 | −0.39 | −0.45 | −0.60 |
Table 1

(Continued)

| ID  | F105W | F110W | F125W | F160W |
|-----|-------|-------|-------|-------|
| 25  | 29.26 ± 0.19 | 29.45 ± 0.30 | 29.08 ± 0.26 | 29.18 ± 0.33 |
| 26  | 29.40 ± 0.19 | 30.03 ± 0.42 | 29.93 ± 0.35 | 29.19 ± 0.14 |
| 27  | 29.75 ± 0.18 | 29.20 ± 0.13 | 29.91 ± 0.31 | 30.51 ± 0.66 |
| 28  | 30.20 ± 0.74 | 28.89 ± 0.16 | ... | ... |
| 29 (bulge) | 26.65 ± 0.02 | 26.65 ± 0.03 | 26.55 ± 0.03 | 26.25 ± 0.04 |
| 30  | 30.41 ± 0.41 | 29.85 ± 0.14 | 29.45 ± 0.19 | 30.40 ± 0.41 |

Notes. The filter throughput is in AB mag system. The listed magnitudes have been corrected for a constant factor (2.26 mag) to remove the magnification effect.

a After k-correction, this filter is used as the rest-frame U band in the comparison with low-z cluster complexes.
b Main spectral features transmitted, [O ii].
c Main spectral features transmitted, Balmer break, [O ii].
d Applied aperture correction; see the main text for details.
e Main spectral features transmitted, Balmer break, [O ii].
f Main spectral features transmitted, Balmer break, [O ii], [Hβ], [O iii].
g Main spectral features transmitted, Hβ, [O iii]. After k-correction, this filter is used as the rest-frame V band in the comparison with low-z cluster complexes.
h Main spectral features transmitted, Hα.

Figure 2. Recovered clump age (left) and mass (right) distributions. The outputs have been obtained for three sets of Yggdrasil models and metallicity $Z = 0.008$: instantaneous burst (red dashed line histogram), continuous SFR for 10 Myr (black solid line histogram, these outcomes will be used in the analysis), continuous SFR for 100 Myr (blue dot-dashed line histogram).

(A color version of this figure is available in the online journal.)

In Figure 2, we show the recovered age and mass distributions of the clumps. We only use the outcomes of the fit obtained with the three SFHs and metallicity $Z = 0.008$ models. We note, however, that the fit with models at lower metallicity ($Z = 0.004$) produces similar results. We preferred the former metallicity assumption because it is closer to the mean value found by Yuan et al. (2011). The SFH of the clumps is unknown. If we look at cluster complexes in the local universe (see the last column in Table 3) their age range varies from a few megayears (in agreement with an assumption of instantaneous burst) to a few tens of megayears (compatible with prolonged star formation episodes). Therefore, we decided to compare the results obtained from the three different SFHs: instantaneous, continuous for 10 Myr (similar to the range observed for more massive and complex star-forming knots), and for 100 Myr. Most likely the real SFH lies between these two last values. In Figure 2, the age distribution recovered with an instantaneous burst model (red dashed line histogram) produces a peak at young ages ($\sim$6–7 Myr) which disappears if a continuous SFR (a more realistic assumption) for a given time model is used instead. The ages we fit are defined from the beginning of each star formation episode of each model. Allowing for a range of SFHs introduces a certain age degeneracy because blue objects can be fitted to older ages for more prolonged star formation scenarios, but since the overall SED is continuously evolving even if the SFR stays constant (e.g., Zackrisson et al. 2005), many of our clumps turn out to be best fitted by ages shorter than the assumed duration of the constant-SFR phase. Models with prolonged SFR have, as expected, the tendency to produce slightly older ages for the clumps. On the other hand, the mass distributions do not change significantly and range between a few times $10^5$ up to $10^8 M_\odot$ (more than 70% have masses below $3 \times 10^7 M_\odot$), in spite of the different models used. The recovered mass range is in good agreement with the results obtained by Elmegreen et al. (2009a). We decided to use the results obtained with the fit to the model of 10 Myr continuous star formation and metallicity $Z = 0.008$ in the rest of the analysis. However, we stress that the conclusions would not change significantly if the results obtained with a 100 Myr SFH history would have been used instead.

3.2. Clump Size

To measure the clump size we used the frame with the deepest exposure, i.e., the F814W. Any attempt to directly measure the FWHM of the Sp 1149 clumps has failed because the S/N is...
not high enough except in a few sources, e.g., 1, 4, and the bulge (29). For these sources we measured FWHMs (2.67, 2.53, 2.61 pixel, respectively) comparable to the mean FWHM of the stellar PSF (σ = 2.64 ± 0.33 pixel).

To determine the sizes we applied an indirect deconvolution of the object PSF with the stellar PSF, which gives an approximate measure of the \( R_{\text{eff}} \) (the half-light radius, i.e., the radius that contains half of the total light of the system, e.g., Portegies Zwart et al. 2010).

We looked at the concentration index (CI) of each clump, e.g., the concentration of the light as a function of the PSF radius. We defined the CI as the difference between the magnitude of the clump at a photometric aperture of 1 pixel and the magnitude at a radius of 3 pixel. In Figure 3, we see the recovered values of CI as a function of their final magnitude in the F814W filter (black dots). CI values close to 0.0 means that all the light of the source is concentrated in the central circle of 1 pixel radius. The stellar PSF CI sets the limit between resolved and unresolved objects in the frame. We used a few stars in the F814W frame to measure this limit, obtaining a CI = 1.15 mag (red cross in Figure 3). Clumps with CI below this limit are considered unresolved.

To trace the behavior of the CI as a function of the size of the source, we estimated the CI of simulated extended objects with increasing \( R_{\text{eff}} \). We constructed for the F814W filter a modeled PSF using the PyRAF task PSF. The stellar PSF was then convolved with an analytic surface profile (Elson et al. 1987; and power index of 3.0) of increasing \( R_{\text{eff}} \) using the task mkcompsf from the publicly available software BAOLab (Larsen 1999). Each extended profile has been used to include several mock sources in the F814W frame, for which we derived a median CI value. We used the following grid of \( R_{\text{eff}} \) values, 0.06, 0.6, 1.3, 2.5, 3.8, 5.1, 6.4, 8.9, and 12.7 pixel (1 pixel = 191 pc at the distance of Sp 1149). The recovered CI values are shown

\[
\begin{array}{cccccc}
\text{ID} & M_U & M_V & R_{\text{eff}} & \text{Age} & E(B - V) \\
\hline
0 & \text{15.52} & \text{15.74} & <120 & \text{3.00} & 1.07 \\
1 & \text{16.20} & \text{15.02} & <120 & \text{1.00} & 0.08 \\
2 & \text{15.91} & \text{14.86} & \ldots & 0.10 & 1.00 \\
3 & \ldots & \text{15.42} & <120 & 6.00 & 1.38 \\
4 & \text{16.15} & \text{14.94} & <120 & 1.00 & 1.00 \\
5 & \text{15.84} & \text{14.96} & <120 & 1.00 & 1.11 \\
6 & \text{16.06} & \text{13.63} & <120 & 0.50 & 0.04 \\
7 & \text{14.93} & \text{15.02} & <120 & 3.00 & 0.54 \\
8 & \text{14.97} & \ldots & <120 & 2.00 & 0.22 \\
9 & \ldots & \text{13.87} & \ldots & 1.00 & 0.06 \\
10 & \text{15.47} & \text{14.79} & \ldots & 0.20 & 0.07 \\
11 & \text{15.57} & \text{15.60} & \ldots & 0.20 & 0.16 \\
12 & \text{15.75} & \text{15.45} & \ldots & 1.00 & 0.48 \\
13 & \text{16.69} & \text{16.70} & \ldots & 1.00 & 14.00 \\
14 & \text{16.26} & \text{16.01} & \ldots & 3.00 & 2.73 \\
15 & \text{15.42} & \text{15.63} & <120 & 0.10 & 0.32 \\
16 & \text{15.83} & \text{15.23} & 244 & 8.00 & 1.23 \\
17 & \text{15.75} & \ldots & \ldots & 1.50 & 0.18 \\
18 & \text{16.31} & \text{16.51} & 976 & 20.05 & 11.80 \\
19 & \text{16.35} & \text{14.42} & 487 & 0.10 & 0.04 \\
20 & \ldots & \text{14.53} & \ldots & 1.50 & 0.07 \\
21 & \text{15.80} & \ldots & \ldots & 0.50 & 0.03 \\
22 & \ldots & \ldots & \ldots & 30.23 & 1.85 \\
23 & \text{16.81} & \text{15.16} & 244 & 3.00 & 2.01 \\
24 & \ldots & \ldots & \ldots & 1.40 & 0.10 \\
25 & \ldots & \text{15.06} & 487 & 3.00 & 0.10 \\
26 & \text{16.20} & \text{14.28} & 244 & 0.10 & 0.05 \\
27 & \ldots & \text{14.28} & 244 & 1.00 & 0.06 \\
28 & \text{16.32} & \text{17.07} & \ldots & 6.00 & 0.91 \\
29 (bulge) & \text{18.37} & \text{17.60} & \ldots & 3.00 & 11.59 \\
30 & \ldots & \text{14.70} & 244 & 3.00 & 0.25 \\
\end{array}
\]

Notes. The rest-frame absolute magnitudes in \( U \) and \( V \) bands have been obtained applying a \( k \)-correction to the F775W and F1255W filters and converted from ABmag to Vega system to allow the comparison with the local reference sample. The ages, masses, and extinction have been derived using models with 10 Myr continuous star formation. See the main text for details.

http://baolab.astroduo.org/
as crosses in Figure 3, in a scale from orange to purple for increasing $R_{\text{eff}}$. The size assigned to each clump corresponds to the $R_{\text{eff}}$ of the simulated CI which is the closest to the observed one. Clumps with CI smaller than the stellar CI (red cross) or photometric errors larger than 1 mag have null measurements. The smallest recovered $R_{\text{eff}}$ we consider significant is $R_{\text{eff}} = 0.6$ pixel $\sim 122$ pc. Sources with CI larger than the stellar one but smaller than half a pixel ($R_{\text{eff}} < 0.5$ pixel) have assigned an upper limit to their measurement corresponding to $\sim 120$ pc. One can see that the clumps for which the FWHM has been directly measured on the image (1, 4, and 29) fall in this category (e.g., sources with an upper limit to the real size). The recovered $R_{\text{eff}}$ are all listed in Table 2. In a recent work, Livermore et al. (2012) detected the bulge of Sp 1149 in a Hα-dedicated survey with HST. They measured a size for our source 29 of 174 $\pm$ 34 pc, which is in good agreement with our determined upper limit.

### Table 3

**Local Reference Sample Properties**

| ID                | $M_U$ (mag) | $M_V$ (mag) | $R_{\text{eff}}$ (pc) | Mass $(10^5 M_\odot)$ | Age$^e$ (Myr) |
|------------------|-------------|-------------|------------------------|------------------------|---------------|
| **Literature data** |             |             |                        |                        |               |
| NGC 6946-1$^b$   | $-15.81$    | $-15.00$    | $390.0$                | $100.00$               | $1$–$30$      |
| Antennae-1       | $-14.85$    | $-13.80$    | $140.0$                | $22.50$                | $20$–$30$     |
| Antennae-2       | $-14.69$    | $-13.60$    | $198.0$                | $8.00$                 | $6$–$16$      |
| Antennae-3       | $-14.69$    | $-13.60$    | $200.0$                | $8.00$                 | $7$–$14$      |
| Antennae-4       | $-15.08$    | $-13.50$    | $136.0$                | $10.00$                | $3.5$–$5.5$   |
| Antennae-5       | $-15.77$    | $-14.00$    | $175.0$                | $11.00$                | $3$–$4$       |
| Antennae-6       | $-14.85$    | $-13.10$    | $90.0$                 | $10.00$                | $1$–$2$       |
| Antennae-1Nuc    | $-16.04$    | $-14.60$    | $123.0$                | $9.00$                 | $4$–$6$       |
| Antennae-2Nuc    | $-16.07$    | $-14.30$    | $140.0$                | $11.00$                | $3.5$         |
| M 51d-C2         | $-13.40$    | $-12.60$    | $160.0$                | $1.40$                 | $7$           |
| M 51-D2          | $-13.20$    | $-12.40$    | $160.0$                | $1.10$                 | $7$           |
| M 51-E2          | $-12.00$    | $-11.20$    | $85.0$                 | $0.40$                 | $7$           |
| M 51-F2          | $-11.90$    | $-11.10$    | $100.0$                | $0.30$                 | $7$           |
| M 51-G2          | $-14.10$    | $-13.30$    | $220.0$                | $2.60$                 | $7$           |
| M 51-B1          | $-12.80$    | $-12.00$    | $125.0$                | $0.70$                 | $7$           |
| **This work**    |             |             |                        |                        |               |
| M 83-0           | $-12.89$    | $-11.56$    | $50.0$                 | $0.44$                 | $1$–$3$       |
| M 83-1           | $-14.12$    | $-12.87$    | $65.0$                 | $1.45$                 | $5$–$10$      |
| M 83-2           | $-14.12$    | $-12.86$    | $70.0$                 | $1.43$                 | $6$           |
| M 83-3           | $-13.55$    | $-12.26$    | $85.0$                 | $0.83$                 | $1$–$13$      |
| M 83-4           | $-13.88$    | $-12.58$    | $110.0$                | $1.12$                 | $1$–$14$      |
| M 83-5           | $-13.86$    | $-12.27$    | $117.5$                | $0.84$                 | $5$–$10$      |
| M 83-6           | $-13.51$    | $-12.23$    | $117.5$                | $0.80$                 | $5$–$10$      |
| M 83-7           | $-13.20$    | $-11.91$    | $80.0$                 | $0.60$                 | $5$           |
| M 83-8           | $-13.74$    | $-12.36$    | $120.0$                | $0.91$                 | $1$–$10$      |
| M 83-9           | $-13.06$    | $-11.69$    | $100.0$                | $0.49$                 | $5$           |
| Antennae-Knot S$^c$ | $-17.25$    | $-16.3$    | $490.0$ (20)           | $40$                   | $6$–$30$      |
| NGC 7252-W3      | $-20.19$ (16.2) | $-19.10$ (15.67) | $17.5$ | $962.70 (800.0)$ | $10$ (400) |
| NGC 7252-W30     | $-18.39$ (14.6) | $-17.30$ (14.07) | $9.3$ | $190.0 (160.0)$  | $10$ (400) |
| Haro 11-Knot C   | $-19.35$    | $-18.34$    | $400.0$                | $420$                  | $3$–$35$      |
| Haro 11-Knot B   | $-18.42$    | $-17.58$    | $400.0$                | $210$                  | $1$–$20$      |
| Haro 11-Knot A   | $-18.79$    | $-17.39$    | $400.0$                | $170$                  | $1$–$20$      |
| ESO 338-KnotA    | $-18.00$    | $-16.67$    | $301.0$                | $27.41$                | $3$–$11$      |

**Notes.** The photometry is in Vega system. The distance modulus used for the new data is 28.25 mag for M 83, 31.71 mag for the Antennae, 34.58 mag for Haro 11, and 32.87 for ESO 338.

$^a$ For the literature data, we report the mean age listed in the corresponding papers. We list, instead, the minimum and maximum age of the clusters located in the complexes of M 83, in Knot S of the Antennae, in the knots of Haro 11 and ESO 338.

$^b$ Larsen et al. (2002).

$^c$ The Antennae data are taken from Bastian et al. (2006a), except the $U$-band luminosity (obtained from the color produced by our models at the corresponding age) and the radius of the complexes 1 and 2 Nuc (corresponding to the aperture radius used to do photometry).

$^d$ The M 51 data are taken from Bastian et al. (2005a), except the $U$-band luminosity (obtained from the color produced by our models at the corresponding age).

$^e$ For the cluster analysis of Knot S and the star formation history, see Whitmore et al. (2010). In this work, we use the size of the whole complex (490 pc). We indicate in brackets the $R_{\text{eff}} = 20$ pc as reported by Bastian et al. (2013). See the main text for details.

$^f$ The observed values at the current age of the clusters from Bastian et al. (2006b) are shown in brackets. This information has been used to extrapolate the luminosity in $U$ and $V$ bands, applying an $M/L$ ratio at 10 Myr given by our models.
of 120 pc, considering that the two measurements have been taken at different wavelengths (Hα is generally more physically extended than U-band rest-frame emission).

4. BUILDING A REFERENCE SAMPLE

In order to compare the properties of Sp 1149 clumps to local and high-z complexes, we assembled a comprehensive reference sample which is used in Section 5.

4.1. Cluster Complexes in Nearby Galaxies

For the local sample we extended the data available in the literature with new contributions. They are all summarized in Table 3.

The cluster complex properties of the Antennae and M 51 are from Bastian et al. (2005a, 2006a, respectively). For two of the Antennae complexes, the sizes were not available, therefore, we assigned the size of the aperture radius used to do photometry. In both studies, only the V-band luminosity was available. We derived the corresponding U-band magnitude from the color information at the given age (mean age from the interval given for each of the Antennae complexes, mean age of 7 Myr for all the complexes of M51, as indicated in the paper). The properties of the massive cluster complex in NGC 6946 are available from Larsen et al. (2002). The sample also includes some of the most massive known intermediate-age super star clusters (SSCs), W3 and W30, located in the post-merger galaxy NGC 7252 (Bastian et al. 2006b). We used the age and $M_V$, information reported in the paper to derive the cluster properties (e.g., $M_U$, $M_V$, mass) at the age of 10 Myr, using SSP models (color and mass-to-light ratio, $M/L_V$).

The new contribution to this local sample comes from cluster complexes in the M 83 spiral, an extended star-forming knot in the Antennae, Knot S (Whitmore et al. 2010), and in two starburst blue compact galaxies Haro 11 and ESO 338-IG04 (hereafter ESO 338). The M 83 photometric data (WFC3/F336W and WFC3/F555W) have been presented in Bastian et al. (2011). The photometry of Knot S has been performed on the latest WFC3/UVIS F555W and F336W acquired images (GO 11577; PI: Whitmore). The reduction of the ACS/F330W and F550M frames, used for the analysis of Haro 11 complexes, is described in Adamo et al. (2010). The description of the data (PC/F336W and ACS/F550M frames) used for the analysis of ESO 338 can be found in Östlin et al. (2009). We did aperture photometry using a radius of 115 pixel (~100 pc), 115 pixel (~490 pc), 40 pixel (~400 pc), and 35 pixel (~300 pc), for M 83, Knot S, Haro 11, and ESO 338, respectively. The fluxes have been corrected for Galactic extinction. No aperture correction has been included. The photometry listed in Table 3 is in the Vega system.

The masses have been estimated using the V-band luminosity and the $M/L$ ratio given by Yggdrasil models (instantaneous burst) corresponding to a stellar population of 5 Myr for cluster complexes in M 83 and the knot of Eso 338, 6 Myr for Knot S, 10 Myr for the knots in Haro 11. We note that these mass estimates have to be considered a lower limit to the real values. For consistency with the data available in the literature we use $M/L$ ratio derived from SSP models at a single age. However, the SFH in these regions is more complex and the $M/L$ ratio is very sensitive to the age of the stellar population. The deviations are important for more complex regions such as the star-forming knots in starburst galaxies. For example, dynamical mass estimates of the knots in Haro 11 give about a factor of 10 higher masses (Guseva et al. 2012).

Because the regions we looked at are resolved, we defined the size of each complex as the radius which included all the bright sources in the clustered region. In cases where the regions are elongated we defined semi-minor/major axes and used in the analysis the mean of the two values. For the two SSCs W3 and W30 we report the values of the $R_{eff}$. However, Bastian et al. (2013) report that their surface profiles are much more extended and observe the same behavior in Knot S. This region is much younger than W3 and W30 but it could be considered a progenitor analog of these SSCs. The $R_{eff}$ measured for Knot S is about 20 times smaller than the radius of the complex (see Table 3). This discrepancy is likely caused by the fact that one of the most massive SSCs observed in the Antennae (Whitmore et al. 2010) is located in the center of Knot S, dominating is light.

Morphologically, we noticed that in spiral galaxies some of the complexes are very diffuse, while others are centrally concentrated and much brighter, but do not extend as far (e.g., see images showed in Bastian et al. 2005a; Whitmore et al. 2010). The cluster complex in NGC 6946 is quite peculiar. Unlike complexes in other spirals it is more extended and massive, hosting a very massive cluster and several other intermediate mass ones. Complexes in blue compact galaxies are more massive, extended and host a large population of clusters, suggesting denser clustering. These regions are real starburst knots (e.g., Adamo et al. 2010) with $\Sigma_{SFR}$ an order of magnitude higher than the values presented in Bastian et al. (2005a, 2006a).

4.2. The High-z Sample

The high-z sample includes contributions from three different observational samples and one simulation work. The first sample includes the properties of the clumps in six SINS (SINFONI Integral field spectroscopy survey of $z \sim 2$ galaxies; Förster Schreiber et al. 2009) targets. NICMOS imaging of these systems has allowed the derivation of the size and mass of

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Figure 3. Observed concentration indices (CI) as a function of their final (de-lensed) F814W magnitude (black dots). On the left side the crosses indicate the reference CI values used to derive the $R_{eff}$ of the clumps. The red cross corresponds to the stellar CI. The other crosses, from orange to purple, show the recovered CI for increasing values of $R_{eff}$ of simulated extended sources. Sources below the red-dashed line are unresolved. The objects which fall between the dashed and dot-dashed lines have upper limits to their sizes. Sources above the dot-dashed line are resolved. The green diamond shows which source is the bulge of the galaxy. See the main text for details.

(A color version of this figure is available in the online journal.)
their clumps (Förster Schreiber et al. 2011). We note that the resolution of the NICMOS data has limited the measurement to clumps with sizes larger than \( \sim 1.4 \) kpc (using the FWHM of the clumps), i.e., an order of magnitude larger than our limit of 120 pc. The stellar masses have been estimated from the luminosity of the clumps in the F160W band (corresponding to the V-band rest frame), assuming the \( M/L \) ratio at the age derived from the \( i_{814W} - F160W \) color. Wisnioski et al. (2012) reported the sizes and masses of the clumps of three targets belonging to the WigleZ Dark Energy Survey. The sizes were measured in the H\( \alpha \) redshifted frame, fitting the light profile of the clumps. The masses are dynamical estimates using the velocity dispersion of the H\( \alpha \) emission line and the assumption of virial equilibrium. The third sample includes sizes (from isophotes) and masses (dynamically) derived from clumps of gravitationally lensed targets at \( z \sim 2–3 \) (Jones et al. 2010). Finally, we also consider the size and mass of simulated clumps formed by violent gravitational instability in high-\( z \) disk galaxies (Ceverino et al. 2012).

5. UNDERSTANDING THE PHYSICAL PROPERTIES OF SP 1149

5.1. Luminosity Properties

In order to compare the luminosity of the Sp 1149 clumps to the brightness of local cluster complexes, we have applied the \( k \)-correction to the photometry in the F775W and F125W filters. We used Yggdrasil models at \( z = 0 \) and the constrained ages to derive the correction factor for each clump, obtaining corresponding absolute magnitude in \( U \) (F336W) and \( V \) (F555W) bands at \( z = 0 \).

The total apparent magnitude of Sp 1149 in the F814W frame is \( I \sim 23.4 \pm 0.3 \) mag (Smith et al. 2009). A comparison with the total luminosity produced by the clumps shows that, in spite of being a prominent feature of Sp 1149, they contribute only 15\% of the total luminosity in the rest-frame \( U \) band. In other words, the bulk of the UV light in the galaxy is produced by a diffuse stellar component. A similar result has already been found by Förster Schreiber et al. (2011) in a sample of \( z \sim 2 \) galaxies. Based on its UV properties, Sp 1149 seems to agree quite well with the trend observed at high \( z \).

In Figure 4, one can see the luminosity distribution of the clumps in Sp 1149 compared to the luminosity distribution of the local reference sample. The luminosity of the clumps is similar to the range of brightness covered by the local sample. To better understand the similarities and differences between the two data sets, we divided the local sample into complexes belonging to spirals (green solid line distribution) and complexes formed in merger starburst systems (purple dashed line). The distribution at \( z \sim 1.5 \) clumps sits in between the two subsamples, about \( 2–3 \) mag brighter than the mean magnitude of the complexes in spirals. One can see that the high-\( z \) distribution partially overlaps with the brightest systems in the Antennae (secondary peak in the purple distribution at lower brightness bins). Definitively, the Sp 1149 clumps are not as luminous as the knots in blue compact galaxies or the most massive single clusters of NGC 7252 (bright-end purple distribution). Only the bulge-like component of Sp 1149 (the brightest source in the galaxy) has a luminosity comparable to those of the star-forming knots. We note that the median value of the Sp 1149 clump luminosity is \( M_U = -15.84 \) mag and \( M_V = -15.02 \) mag. These median values overlap quite well with the observed magnitudes of the NGC 6946 massive cluster complex (listed in Table 3), one of the most massive studied complexes in local spirals.

The trend observed in the luminosity function is better understood in a color–magnitude diagram (Figure 5), in which the information coming from the two filters can be explored and correlated. In general, we observe that the high-\( z \) clumps are distributed in a wide range of \( U - V \) color. Clumps with \( U - V < -0.5 \) mag overlap quite well with the range of \( U - V \) color of the local cluster complexes, suggesting that they have similar young ages. These young clumps seem to be the ones with lower stellar mass and are comparable to the complexes in the Antennae or in NGC 6946. On the other hand, clumps with redder colors, e.g., \( U - V > -0.5 \) mag seem to be older and more massive than their younger counterparts. However, because of the extinction, some of these redder clumps could still be young. Using the derived ages we see that about 60\% of the clumps have an age between \( 1–30 \) Myr (corresponding to \( U - V = -0.5 \)). The remaining fraction is definitively older. The coexistence of two populations of clumps (young, low mass, and less bright versus old, more massive, and bright) in Sp 1149 is clearly outlined by the distribution of the clumps in the age–mass and age–\( M_V \) diagrams shown in Figure 6. Of the local reference sample, only the two very massive clusters in NGC 7252 (see position indicated by the arrows) have observed colors and luminosities comparable to the older clumps. Their location in the diagram also suggests that of the local complexes;
the knots in blue compact galaxies and the Antennae seem to be good progenitor analogs of these evolved clumps.

5.2. Simulating Local Galaxies at the Redshift of Sp 1149

From the comparison of the luminosity properties, we observe that none of the local cluster complexes in our reference sample have ages (colors) comparable to the more evolved clumps of Sp 1149. This evidence suggests that complexes in local galaxies are destroyed faster than in high-z galaxies. The reason for such short lifetimes could be related to the low masses of the complexes compared to those of the clumps.

In local galaxies, star formation is observed to happen in hierarchical regions, and, in general, giant molecular cloud (GMC) complexes show sizes up to ~1 kpc (e.g., Elmegreen 2011b). Studies of the clustering in local star-forming galaxies show tight spatial and temporal correlations between massive stars and clusters at young ages (e.g., Bastian et al. 2007). However, these hierarchical structures are rapidly erased with time.

To test how local star-forming complexes would look at high z we have simulated images of local galaxies (M 83, Antennae, Haro 11) to have the same depth and resolution as observed for Sp 1149 (we remind the reader that the linear physical scale of the latter has been magnified by $\sqrt{8} \times$ and the fluxes by a factor of $8 \times$). Following the method presented in Overzier et al. (2008), we have first resampled each galaxy science frame to a pixel size that matches the angular resolution of the Sp 1149 data (both smoothing and rebinning the frames by the corresponding quantity). As a second step, we have taken into account the dimming of the surface brightness as a function of the luminosity distance. The final rescaled images are included in Figure 1. We simulated two filters for each galaxy, $U$ and $V$ bands. One can clearly see that at the distance and resolution of Sp 1149 the star-forming regions in local galaxies blend in clumps, similar to what we observe in our high-z target. The brightest clumps are, however, all located in correspondence with the youngest star-forming regions in the galaxies. For example, in the case of M 83, a comparison between the location of the bright clumps in the high-z $U$-band image and the contours of pronounced Har regions presented in Figure 3 by Silva-Villa & Larsen (2012) shows that they nicely trace each other. A similar trend is observed in the Antennae (e.g., compare the location of the clumps with the color-composite image in Figure 1 by Whitmore et al. 2010) and in Haro 11 (e.g., Figure 1, Östlin et al. 2009).

Our simple exercise suggests that, indeed, the bright clumps in our local reference sample are very young (emitting Hα) or still star-forming and, if present at all, older clumps would clearly be less luminous. Therefore, in spite of the observational biases, the high-z Sp 1149 clumps appear to be in some aspects intrinsically different. The young clumps of Sp 1149 resemble quite closely the cluster complexes observed in local galaxies. Their difference in mass (luminosity) could be caused by the poor resolution at high z. The older clumps of Sp 1149 are more massive and long-lived than any of the local complexes in our sample. The star-forming knots of blue compact galaxies

![Figure 5](image-url)  
Figure 5. Color–magnitude diagram of the clumps (black dots) and the local reference sample (symbols are labeled in the inset). The black arrow shows in which direction of the diagram the data points would move after an extinction correction is applied. Photometric error bars are included. The errors of the local sample are very small (order of 0.005 mag) so are not shown. The green diamond enclosing the brightest black dot shows the position of the bulge-like source. The orange stars indicate the position of the two very massive star clusters of NGC 7252 at the age of 10 Myr. The orange arrows show the current position, corresponding to the age of 400 Myr. The overplotted tracks are from Yggdrasil models at $z = 0$, assuming $Z = 0.008$, covering factor of one, Kroupa IMF, and 10 Myr of constant SFR. Different stellar masses are assumed (labeled in the plot). A square and a triangle symbol indicate the starting and ending of the tracks, while the asterisk shows the position in the track corresponding to 10 Myr.

(A color version of this figure is available in the online journal.)
are good progenitor analog candidates of these evolved clumps. However, to test this scenario one should also be able to detect such compact knots in local post-starburst environments and to date none have been found. Finally, the two very massive clusters of NGC 7252 have luminosity and color similar to the old clumps, showing that also at low redshift it is still possible, under extreme environmental conditions (e.g., a galaxy merger), to form very massive systems. However, in the following section we will see that these clusters are structurally quite different from the Sp 1149 clumps.

5.3. Mass versus Size Relation

In M 51, Bastian et al. (2005a) found that the stellar mass of the cluster complexes scales with their size, i.e., \( M \propto R^2 \), in a similar way as for the GMCs of the Milky Way and nearby galaxies (Solomon et al. 1987; Bolatto et al. 2008). This relation, assuming that the GMCs are in virial equilibrium, suggests that they have constant surface gas density, i.e., \( \Sigma \approx 100–200 M_\odot \text{ pc}^{-2} \). Therefore, at least in nearby galaxies, GMCs appear to form under similar physical conditions. Interestingly, the mass–size relation does not hold for young star clusters (Bastian et al. 2005b), formed in the denser regions of the GMCs. Because of turbulence acting as a scale-free process, collapsing GMCs fragment in dense clumps and cores. As such, not the single star clusters but the whole complex which formed out of the GMC gas reservoir (assuming a certain star formation efficiency; e.g., Bastian et al. 2005a), can be considered as tracers.

In this context, thanks to the resolution achieved in Sp 1149, we can try to understand the physical conditions that have led to the formation of its clumps. Using the constrained physical properties, we attempt to build a mass–size plot, where we also include the properties of the local cluster complexes and of the high-\( z \) sample. The plot is shown in Figure 7. Although the whole data set is included in the plot, we point out that the two SSCs of NGC 7252 have been excluded from the fit because as pointed out above star clusters do not follow the mass–size relation. Similarly, the data obtained from Ceverino et al. (2012) have been included only for illustrative purposes. Finally, the Sp 1149 clumps for which we derive an upper limit to their sizes are also included. The Sp 1149 clumps for which we derive an upper limit to their sizes have been excluded from the fitting analysis. We first attempted a single fit to the selected sample. We recovered a slope much steeper than expected from the mass–size relation. Similarly, the data obtained from Ceverino et al. (2012) have also been excluded from all the fits. The lower right insets show the recovered slopes from the fit of different samples. The \( \alpha_{M/L} \) has been obtained fitting all the clumps and complexes, in spite of their redshift. The \( \alpha_{\text{loc+Sp1149}} \) is the slope of the sample including only the local reference sample and the clumps of Sp 1149. On the other hand, the \( \alpha_{\text{high-z}} \) has been obtained fitting only the high-\( z \) sample, excluding the Sp 1149 clumps.

(A color version of this figure is available in the online journal.)

In Figure 8, we compare the mass–size relation obtained for the cluster complex sample and the clumps of Sp 1149 with a local sample of GMCs. As already noticed by Bastian et al. (2005a), the vertical offset between the GMC sample and cluster complexes in local spirals could be due to the star formation efficiency (only a fraction of the gas in the GMCs fragment in dense clumps and cores. As such, not the single star clusters but the whole complex which formed out of the GMC gas reservoir (assuming a certain star formation efficiency; e.g., Bastian et al. 2005a), can be considered as tracers. In a recent work, Livermore et al. (2012) showed evidence of a relation between the size and the SFR of the clumps (or their \( H\alpha \) surface brightness). They observed a systematic offset toward higher surface brightness at higher redshift. This trend could be understood in a scenario where high-\( z \) galaxies have disks with higher gas surface density able to form more massive clumps. Indeed, we observe that the galactic SFRs of our high-\( z \) sample, which can be considered an indirect measurement of the total gas surface density of the galaxy (e.g., Bigiel et al. 2008), are between a factor of 2 and 20 higher than in Sp 1149. Therefore, the vertical shift of the high-\( z \) sample in the mass–size relation could be evidence of different conditions under which those clumps have formed (e.g., Swinbank et al. 2012). However, we stress that this offset may rise because of the different methods used to derive the masses.
GMCs will form stars). A considerable fraction of complexes of merger systems and of Sp 1149 (including a large fraction of clumps with upper limits on the size) are offset with respect to complexes in local spirals. This trend could be evidence that these more massive systems have formed in GMCs with higher gas surface density (i.e., more massive). The bulge-like clump shown as a red arrow deviates clearly from the relation, suggesting a different mechanism for its formation as discussed in the following section.

6. CLUMP FORMATION AND SURVIVAL AT HIGH z

As discussed in Section 1, clumps at high z have likely formed in marginally unstable disks in large gravitational instabilities necessary to overcome the high turbulence of the gas. In the local universe, only merging systems, e.g., the Antennae, appear to be able to form very massive GMCs (e.g., Wei et al. 2012) between $10^4$ and $10^5 M_\odot$. The formation of such massive GMCs could be related to the enormous external pressures that the merger event causes on the gas. Star-forming galaxies like the Milky Way form GMCs via agglomeration reaching maximum masses $< 10^7 M_\odot$ (Williams & McKee 1997). In this context, we try to understand how the clumps in Sp 1149 have formed.

First of all, we focus on the young population of clumps (age $< 30$ Myr). These clumps have formed during the recent star-forming events, and are likely related to the current SFR. Smith et al. (2009) reported a current SFR for Sp 1149 of $\sim 6 M_\odot$ yr$^{-1}$ from the $U$-band rest-frame fluxes while using the total Hβ luminosity, Livermore et al. (2012) found a slightly smaller value, i.e., $\sim 1.5 M_\odot$ yr$^{-1}$. Thus, the SFR of Sp 1149 is not very different from the SFR in local spirals. Using the optical rest frame we derive a crude approximation of the SFR in these more massive and evolved clumps. We obtain that for ages between 30 and 50 Myr, the averaged clump SFR $\sim 3 M_\odot$ yr$^{-1}$; in the range 50–100 Myr $\sim 3 M_\odot$ yr$^{-1}$; and $\sim 22 M_\odot$ yr$^{-1}$ in clumps with ages older than 100 Myr. If we consider that each clump contributes to a fraction of the total SFR in the host system, then this increasing value of the SFR as a function of time may indicate a change in the conditions under which clumps have formed in the past (higher surface gas density, larger gravitational instabilities) and are forming (see above, e.g., similar conditions to local star-forming galaxies) in Sp 1149. However, we stress that this is only a suggestion and that it depends on the many assumptions we have made in deriving the clump properties and the initial stellar masses.

6.1. Bulge Formation in Sp 1149

In high-z galaxies, the formation of the bulge seems to be a gradual process. Elmegreen et al. (2009a) observed that a considerable fraction of clump clusters and chains have red and massive clumps which could be considered bulge progenitors. These young bulges are a few times more massive than the other star-forming clumps in the system but the ages are similar.
In young spirals, bulges are older and about 20–30 times more massive than the clump population, suggesting a clear evolution. Simulations (e.g., Bournaud et al. 2008; Ceverino et al. 2010) show that clumps migrate toward the center of the system because of dynamical friction and tidal torques in about 10 dynamical times, which corresponds to a timescale of about 0.5–1 Gyr. There, clumps coalesce and slowly build up the bulge and the thick disk component (Bournaud et al. 2007; Elmegreen et al. 2008). The migration of clumps should cause a significant age spread, with older clumps preferentially closer to the center. Indeed, evidence of such a trend has already been reported in previous clump studies (e.g., Elmegreen et al. 2009a; Förster Schreiber et al. 2011).

In this context, we try to use the recovered clump properties of Sp 1149 to test the bulge formation scenario. Source 29 is most likely the young bulge of Sp 1149. Smith et al. (2009) already observed that the bulge does not dominate the luminosity of the galaxy, suggesting its early formation stage. The fit to its SED gives a best age estimate of 30 Myr and a mass of a few $10^8 M_\odot$, in good agreement with the age and mass range reported by Elmegreen et al. (2009a) for young bulges in spirals at $z \sim 1.5$. The age and mass determined by using the models with 100 Myr continuous SFR do not change significantly (age of $40^{+39}_{-39}$ Myr, mass of $6^{+15}_{-5.1} \times 10^7 M_\odot$) our conclusions. The fact that the progenitor bulge is younger than other clumps in the galaxy is not in contrast with the scenario of bulge formation presented.
above. Indeed, star formation can proceed for a longer time in the central region than in single clumps without preventing the accretion of older systems (Elmegreen 2011a). The youth of the bulge is also supported by narrowband H α spectroscopy of the galaxy, where it appears to be quite a bright H α source (Livermore et al. 2012; Yuan et al. 2011).

In Figure 9, we show the position of the detected clumps in the galaxy (top panel), the projected distances (bottom left), and the total dynamical time (bottom right) that the clumps have survived \( (r = age/t_{\text{dyn}}) \) in the galaxy as a function of their age. The oldest clumps (age > 50 Myr) are all at a distance less than 4 kpc. Defining the dynamical time as the ratio between the rotation velocity and the radius, \( t_{\text{dyn}} \propto r/u_{\text{rot}} \), we can get a rough estimate of the possible migration of the clumps (e.g., Förster Schreiber et al. 2011). Yuan et al. (2011) found a rotation velocity for Sp 1149 of \( \sim 150 \) km s\(^{-1}\). For the two old clumps at 1.5 kpc distance, we estimate a \( t_{\text{dyn}} \sim 10 \) Myr. This implies that they have survived about 20 \( t_{\text{dyn}} \), and most likely have migrated to their current position closer to the center. For the remaining old clumps, we see that they have ages about 5–15 times \( t_{\text{dyn}} \), still supporting the scenario of long-lived migrating clumps.

The remaining younger clumps have maximum ages of a few \( t_{\text{dyn}} \), suggesting that their current position is close to where they formed, as expected for young cluster complexes in nearby star-forming galaxies. We note that four old clumps out of six for which we have size measurements are among the largest systems (see black dots encompassed by red squares in Figure 7), suggesting possibly an ongoing dissolution of these systems while approaching the center.

6.2. Formation of Globular Cluster Progenitors

In a recent work, Shapiro et al. (2010) suggested that star clusters forming in clumps at \( z \sim 2 \) may form the metal-rich globular cluster population which is usually associated with the thick disk and bulge component of late spiral galaxies. We closely follow their method, adding, however, some extra information on the cluster formation efficiency of each clump, i.e., the fraction of star formation happening in clusters per clump. There is observational evidence that this fraction (referred to as \( \Gamma \); Bastian 2008) of clusters that a galaxy may form scales with the global SFR of the galaxy (e.g., Goddard et al. 2010; Adamo et al. 2011). Analytical models suggest that such a trend is naturally reproduced if the hierarchical mass spectrum of the interstellar medium is combined with the disk stability (Q parameter) and the limit imposed by the environment on the dissolution of clusters (Kruijssen 2012). For the younger clumps (age < 30 Myr), we use the current Galactic \( \Sigma_{\text{SFR}} \) and the relation produced by Goddard et al. (2010) to derive \( \Gamma \approx 14\% \). Therefore, we will assume that only 14\% of the estimated stellar mass of the clumps will form bound clusters. For the older clumps we do not have direct measurements of the Galactic SFR at the moment they formed, but we have seen in the previous section that blue compact galaxy knots could be possible progenitor analogs of these clumps. Adamo et al. (2011) found \( \Gamma \approx 40\%–50\% \) in these systems, so we assume that 50\% of the initial total stellar mass is in globular cluster progenitors. We approximate the cluster mass function to be a power law of the form, \( dN \propto M^{-2}dM \). The minimum and maximum cluster mass that a clump can form are fixed to \( 10^4 \) and \( (M_{\text{lump}} \times \Gamma) \), respectively. The latter assumption means that the maximum mass of a cluster can be as large as the total fraction of stellar mass of the clump going into clusters. After integrating, we obtain that \( \sim 300 \) star clusters with masses larger than \( 10^4 M_\odot \) have been formed in the young clumps.

In the older clumps (age > 30 Myr) we estimate that about 31,000 young star clusters have been produced. However, as discussed in Shapiro et al. (2010), many of the low-mass clusters will dissolve because of two-body relaxation, while the high-mass end of the power law is better described by a Schechter function with a characteristic high-mass cutoff. Taking these factors into account they estimate that for every 7000 star clusters only a small fraction (\( \sim 120 \)) will survive until \( z \approx 0 \) as globular clusters. Therefore, we predict a lower limit number of about 540 metal-rich globular clusters in Sp 1149 in good agreement with expectations (Shapiro et al. predicted a metal-rich globular cluster population up to \( \sim 600 \) systems for \( z \sim 0 \) galaxies with stellar mass of a few times \( 10^{11} M_\odot \)).

7. CONCLUSIONS

Our single-object study has revealed some useful insights into the intermediate evolutionary stage in which clumpy turbulent systems become modern spiral galaxies.

Sp 1149 is not a special galaxy. Its current SFR (1–6 \( M_\odot \) yr\(^{-1}\)) is similar to what is observed in local spirals and the clumpy appearance is not much different than in other spirals observed at similar redshift. What makes Sp 1149 an interesting case study is its position behind a massive galaxy cluster. The magnification gained by the gravitational lens has allowed us to study the properties of the clumps at a spatial resolution \( \sim 3 \) times higher than other systems at a similar redshift. Thanks to the achieved resolution, we have analyzed what we believe to be the transient evolutionary phase that high-\( z \), intense star-forming galaxies experience when they evolve into more quiescent massive systems, as we observe in the local universe.

Using the CLASH survey, we detected roughly 30 clumps in Sp 1149, for which we derived ages, masses, extinctions, and effective radii. To allow a comparison with local star-forming galaxies, we collected available literature data of cluster complexes with properties of star-forming knots in luminous blue compact galaxies and some of the most massive studied intermediate-age SSCs.

We found that the clumps in Sp 1149 are typically more luminous than cluster complexes in nearby spirals, similar to the cluster complex luminosity of the Antennae, but not as bright as star-forming knots in blue compact galaxies. Curiously enough, the median rest-frame \( U \)- and \( V \)-band magnitudes of the clumps are quite similar to the massive cluster complex in the nearby spiral NGC 6946, which gives a close idea of the morphological aspect and physical properties of the young clumps. Diagnostic tools like the color–magnitude and the age–mass diagrams show the coexistence of two populations of clumps in Sp 1149, with an age spread up to a few hundreds of megayear.

The older clumps are the most massive and luminous. None of the nearby cluster complexes has color (ages) comparable to this subsample of old clumps. This suggests that cluster complexes dissolve faster in contrast to older clumps which have formed in more extreme environmental conditions and have masses (properties) that ensure their survival for a significant amount of galactic dynamical timescales. If we extrapolate the luminosity of these older clumps at the moment they formed they must have been at least 2 mag brighter than currently observed. We suggest that among the nearby reference sample, knots in blue compact galaxies could be considered progenitor analogs of these evolved population.

The younger clumps (age < 30 Myr) have colors that overlap quite well with the local sample of cluster complexes, supporting their young age estimate. The masses are a factor of 10 larger
than typical cluster complexes in local spirals and comparable
to complexes in the Antennae and NGC 6946. In our imaging
simulation of local galaxies at the distance and resolution of
Sp 1149, we observed that indeed the cluster complexes are all
distributed along the bright He regions, i.e., they have recently
formed. The morphological detail of the single complexes are
erased, and we do not see any major difference between them
and the young clumps in Sp 1149, suggesting that the difference
in mass and luminosity could be caused by the limited resolution
we can achieve in the latter system.

We tried to understand how clumps have formed, looking at
the mass–size relation, originally constrained on GMCs of the
Milky Way by Solomon et al. (1987), but similarly reproduced
by cluster complexes in M 51 (Bastian et al. 2005a). In its
original formulation this relation suggests that GMCs have
follow the trend of the mass–size relation. Cluster complexes still follow
the relation because they are scaled down (in total mass)
similar gas surface densities. Cluster complexes still follow
the original formulation this relation suggests that GMCs have
follow the trend of the mass–size relation.

Given the current Σ SFR of the galaxy and using the Kennicutt
law we derive a galactic surface gas density of ~50 $M_\odot$ pc$^{-2}$,
similar to M 83 and M 51, where GMCs can reach masses of
~$10^4$–$10^5$ $M_\odot$. These are most likely the GMC progenitors from
which the young clumps of Sp 1149 have formed, assuming a
star formation efficiency of 10%–50%. However, we find
evidence that the older clumps have formed in more massive
GMCs. We tried to reconstruct the initial conditions for these
evolved systems, assuming that they have lost 50% of their
initial stellar mass because of tidal interaction with the diffuse
stellar population of the disk and about 20% of their initial stellar
mass because of stellar evolution. We obtained that the GMCs
from which these clumps formed should have been at least a
factor of 2.4 more massive than their current derived properties.
We derived that the SFR in the old clumps, at the moment they formed,
was in the range 3–22 $M_\odot$ yr$^{-1}$, suggesting that the total
SFR in the galaxy should have been significantly higher than what is measured currently. The galaxy is experiencing a slow
decline in SFR and is experiencing a likely transitional phase
toward more quiescent star formation modes.

The presence of long-lived clumps allowed us to test another
important element in the evolution of young spirals into a more
evolved system, i.e., the formation of the bulge. Sp 1149 already
has a young bulge, suggesting that the formation process is
ongoing. We found that clumps older than 50 Myr are all
located at a projected distance smaller than 4 kpc and have
survived between 6 and 20 galactic dynamical times. Their
current location in the galaxy must be quite different from their
birthplace, producing evidence of an inward migration toward
the center, where they are expected to build up the bulge. On
the other hand, the dynamical timescales of the younger clumps are
significantly shorter, meaning that they are still close to or
quite close to their birthplace, i.e., along the spiral arms, similar
to the cluster complexes observed in nearby systems.

Finally, we attempted to estimate whether the clumps we
observe in Sp 1149 could produce the expected metal-rich
globular cluster population usually associated with the bulge
and thick disk components of local spirals. We found a consistent
fraction of star clusters that will survive and eventually evolve
into globular clusters. We note that the large majority of
clusters that will survive are produced in the long-lived, more
massive clumps. A large fraction of clusters that are formed in
complexes younger than 30 Myr most likely will not survive long
enough. The difference could be understood in terms of size-of-
sample effects. More massive clumps have a higher probability
than cluster complexes to sample the cluster mass function at
the high-mass bins. These more massive clusters have higher
probability to survive up to a Hubble time.

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