ANALYSIS OF SMALL OSCILLATIONS OF COMPLEX ELECTRICAL SYSTEMS

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Abstract. The article discusses the results of the analysis of the static stability of complex electrical systems. The efficiency of the combined application of the equations of nodal voltages (ENV) and the Lyapunov function in quadratic form for the analysis of small oscillations of the electrical system is shown in the literature, which is called the Allaev method. A joint solution of the equations of nodal voltages and the matrix Lyapunov equation is given, which makes it possible to determine the stability conditions for the electrical system and identify the generator first approaching the stability limit. A study of small oscillations of complex electrical systems, which can be performed in full on the basis of matrix methods, successfully developed in recent decades, is carried out, which is associated with a sharp increase in the speed of computation and the amount of memory of modern computers.

Introduction

The highest danger that disrupts normal power supply is an emergency mode in the electric power system (EPS), called a system accident [9].

To prevent such an emergency mode, it is necessary to constantly assess the static stability of the EPS or its resistance to "small" fluctuations, since it is the violations of such stability that lead to systemic accidents with their negative consequences [11].

The problem of studying the static stability of modern complex electrical systems becomes more complicated due to the presence of weak links in their compositions, the presence of various control devices that prevent the determination of the general setting, units with different time constants, etc. [12]. At present, to solve practical problems, methods are used based on calculating the synchronizing power of one of the power plants of the system, i.e. determining the aperiodic stability under the assumption of the absence of self-swinging [9, 10].

The computational analysis of the static stability of EPS of varying complexity shows that the most rigorous theoretically, convenient for calculations and efficient in terms of results is the use of two fundamental methods - the method of Lyapunov functions in quadratic form and the equations of nodal voltages (ENV), which is called the Allaev method in the literature [1, 4]. Moreover, the methods for studying small oscillations taking into account self-swinging are complex, therefore, sufficiently reliable results can be obtained with a rigorous mathematical description of the control system for controlled objects using their parameters and characteristics based on the matrix approach [7, 8].

1.1 Formulation of the problem

It is known [1-4] that the Lyapunov function in quadratic form for linear differential equations is the only one that provides both necessary and sufficient conditions for the stability of the system under study, when small disturbances arise in it. Therefore, the basis of research in this work is the Lyapunov function in quadratic form and nodal equations, and the subject of research is the linearized differential equations of EPS elements. The matrix equations of the elements of electrical systems, which are the main part of the EPS, are based on the equations of state variables, which are most widely used. The considered matrix equations are used to analyze transient processes and static stability of EPS and to synthesize the optimal parameters of the regulators of synchronous machines operating in a complex electrical system.

The steady-state mode of the studied EPS is determined based on the equations of nodal voltages. The nodal equations, which have a functional relationship between the currents and voltages of the nodes, most fully describe the electrical state of the network of any complexity [5]. In the general case, the nodal equations can be written in the form [1, 5, 7]:

\[ YU = I + Y_0 U_0 + J, \]

where

\[ Y = \begin{bmatrix} y_{11} & -y_{12} & \cdots & -y_{1n} \\ -y_{21} & y_{22} & \cdots & -y_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ -y_{n1} & -y_{n2} & \cdots & y_{nn} \end{bmatrix} \] (1)

and

\[ J = \begin{bmatrix} j_{11} \\ j_{21} \\ \vdots \\ j_{n1} \end{bmatrix} \] (2)
conductivity matrix of the system under study; I, Y, J - matrices, respectively, of the nodal currents, the conductances of the connection with the balancing node, the current sources, which are transverse branches with the given conductances [13].

To solve the nodal equations, we will choose Newton's method in polar coordinates, the advantages of which include the quadratic convergence of iterative processes, the possibility of further use for solving optimization problems and calculating stability [5]. In addition, the node voltages \( U \) and the load angles of the generators \( \delta \), which are used in the Lyapunov equations, determined on the basis of solving the nodal equations, contain all information about the state of the system, no matter how complex it is [6]. On the basis of the calculated values of the voltages of the generating units and units containing rotating machines, the positiveness of the matrices of Lyapunov's quadratic forms is sequentially checked. In essence, the task of analyzing the static stability of a system at small oscillations (aperiodic violation, self-swinging) are contained in \( q \) other variables. All types of violation of the stability of the electrical system are satisfied for an EES of arbitrary complexity [1, 7].

### 1.2 Lyapunov function method in quadratic form

In the classical case, the equations describing the processes in the EPS are homogeneous linear (linearized) differential equations and have the form [1, 2, 4, 10]:

\[
\begin{align*}
\frac{dx_1}{dt} &= a_{11}x_1 + a_{12}x_2 + \cdots + a_{1n}x_n; \\
\frac{dx_2}{dt} &= a_{21}x_1 + a_{22}x_2 + \cdots + a_{2n}x_n; \\
\frac{dx_n}{dt} &= a_{n1}x_1 + a_{n2}x_2 + \cdots + a_{nn}x_n;
\end{align*}
\]

either in matrix form:

\[ \mathbf{\dot{x}} = \mathbf{A} \mathbf{x}, \]

where

\[ \mathbf{A} = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1n} \\ a_{21} & a_{22} & \cdots & a_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nn} \end{bmatrix} \]

and \( \mathbf{x}^T = [x_1, x_2, \ldots, x_n]^T \)- transposed vector of state variables.

To determine stability, use the Lyapunov method and define a function in the form of a positive definite quadratic form

\[ V(x) = \mathbf{x}^T \mathbf{Q} \mathbf{x}, \]

or \( V = \sum_{i,j=1}^n q_{ij}x_i x_j \). The derivative of this function:

\[
\frac{dV}{dt} = \frac{d(\mathbf{x}^T \mathbf{Q} \mathbf{x})}{dt} = \left( \frac{dx}{dt} \right)^T \mathbf{Q} \mathbf{x} + \mathbf{x}^T \mathbf{Q} \frac{dx}{dt} = \\
= (\mathbf{A} \mathbf{x})^T \mathbf{Q} \mathbf{x} + \mathbf{x}^T \mathbf{Q} \mathbf{A} \mathbf{x} = \mathbf{x}^T (\mathbf{A}^T \mathbf{Q} + \mathbf{Q} \mathbf{A}) \mathbf{x}.
\]

Require the Lyapunov function to satisfy the requirement

\[ \frac{dV}{dt} = -W, \]

where \( W = \mathbf{x}^T \mathbf{C} \mathbf{x} \) - arbitrary positive definite symmetric matrix.

Equating expressions (7) and (8), we obtain the equation:

\[ \mathbf{A}^T \mathbf{Q} + \mathbf{Q} \mathbf{A} = -\mathbf{C}. \]

Equation (9), called the matrix Lyapunov equation, provides the stability condition if the inequalities \( \mathbf{V} \geq 0 \) and \( \Psi < 0 \) are simultaneously satisfied in some domain of the space of variables \( (x_1, x_2, \ldots, x_n) \), including the origin [4]. Note that both matrices \( \mathbf{Q} \) and \( \mathbf{C} \) are symmetric. Indeed, if the matrix \( \mathbf{Q} \) is symmetric, that is, \( \mathbf{Q}^T = \mathbf{Q} \), then

\[ \mathbf{C}^T = (\mathbf{A}^T \mathbf{Q} + \mathbf{Q} \mathbf{A})^T = -\mathbf{Q}^T \mathbf{A} - \mathbf{A}^T \mathbf{Q} = -\mathbf{C}. \]

and hence the matrix \( \mathbf{C} \) is symmetric. Since the matrix \( \mathbf{Q} \) is symmetric, the Lyapunov equation is equivalent to the system of \( n(n+1)/2 \) linear algebraic equations [4].

According to Sylvester's theorem [4, 9], the positivity of the principal diagonal minors of the coefficient matrix \( \mathbf{Q} \) of the quadratic form (6) is a necessary and sufficient condition for the stability of the system under consideration under small perturbations. For example, for a simple EPS circuit, it looks like this:

\[ \mathbf{Q} = \begin{bmatrix} q_{11} & q_{12} & q_{13} \\ q_{21} & q_{22} & q_{23} \\ q_{31} & q_{32} & q_{33} \end{bmatrix} > 0, \]

i.e., \( \Delta_{11} = q_{11} > 0 \), \( \Delta_{22} = \begin{vmatrix} q_{11} & q_{12} \\ q_{21} & q_{22} \end{vmatrix} > 0, \)

\[ \Delta_{33} = \begin{vmatrix} q_{11} & q_{12} & q_{13} \\ q_{21} & q_{22} & q_{23} \\ q_{31} & q_{32} & q_{33} \end{vmatrix} > 0. \]

Analysis of violation of the condition \( \Delta_{11}, \Delta_{22} > 0 \) shows that all types of violation of the stability of the electrical system at small oscillations (aperiodic violation, self-excitation, self-swinging) are contained in \( q_{11} \). All other minors of \( \mathbf{Q} \) are positive if \( q_{11} > 0 \). This condition is satisfied for an EES of arbitrary complexity [1, 7].

### 1.3 Mathematical model of the electric power system

The linearized equations of the simplest EPS in the presence of automatic excitation regulators (AER) of proportional or strong action on the synchronous generator have the form [2, 4, 10]:

- equation of the relative motion of the rotor of a synchronous machine:

\[ T_{d}(d^2\Delta/dt) = -P_{d}(d\Delta/dt) - \Delta P; \]

- transient equation in the excitation winding:

\[ T_{\phi}(d\Delta_{\phi}/dt) = \Delta E_{\phi} - \Delta E_{\phi}; \]

- transient equation in the field winding:

\[ T_{A}(d\Delta_{\phi}/dt) = \Delta \phi - \Delta E_{\phi}; \]

- converting element equation:

\[ \psi = \Delta \phi; \]

- measuring element equation:

\[ \psi = \Delta \phi; \]
1.4 Numerical results

Consider the application of the equations of the nodal voltage and the Lyapunov function in quadratic form using the example of a three-node circuit (Fig. 1). As a balancing one, we will choose a zero node, the first and third nodes are generating, the second node is load.

![Fig. 1. Schematic diagram of a three-node electrical system](image)

The analysis of the static stability of a complex EPS will be carried out on the basis of known assumptions [1, 12]:
- when calculating the synchronizing power of any of the generators, the rotor angles of all other generators remain unchanged;
- the system of differential equations of the relative motion of the rotors of synchronous generators are described in the form

\[
\begin{align*}
T_\alpha \frac{d^2 \delta_1}{dt^2} + P_\alpha \frac{d\delta_1}{dt} &= P_\alpha - P_r(\delta_{12}, \delta_{13}, \ldots \delta_{1n}) ; \\
T_\beta \frac{d^2 \delta_2}{dt^2} + P_\beta \frac{d\delta_2}{dt} &= P_\beta - P_r(\delta_{12}, \delta_{13}, \ldots \delta_{1n}) ; \\
T_\gamma \frac{d^2 \delta_3}{dt^2} + P_\gamma \frac{d\delta_3}{dt} &= P_\gamma - P_r(\delta_{12}, \delta_{13}, \ldots \delta_{1n}) ;
\end{align*}
\]

- power equations for synchronous generators, expressed in terms of the intrinsic and mutual conductances of the equivalent circuit branches:

\[
P_1 = E_1^2 y_{11} \sin(\delta_1 - \alpha_1) + \ldots + E_1 y_{1n} \sin(\delta_1 - \alpha_n) + E_n y_{1n} \sin(\delta_n - \alpha_1) + \ldots + E_{2n} y_{2n} \sin(\delta_n - \alpha_n) ;
\]

\[
P_n = E_n^2 y_{mn} \sin(\delta_n - \alpha_m) + \sum_{i=j} E_i y_i \sin(\delta_i - \alpha_n),
\]

where \( \delta_1 \) and \( \delta_n \) - absolute and relative load angles of generators; \( E_i \) - generator electromotive forces; \( T_\alpha \) - constant inertia of aggregates; \( P_\alpha \) - generator equivalent damping factors; \( P_r \) - electromagnetic powers of synchronous generators; \( y_{mn} \) - own and mutual conductivities of the system; \( \alpha_m \) and \( \alpha_n \) - corresponding padding angles.

Below are the initial data and parameters of a complex electrical system (Fig. 1).

Parameters of nodes:
- G1: \( P_{G1}=300 \text{ MW}; \cos \varphi_{G1}=0.8; U_{G1}=500 \text{ kV}; T_{j1}=6 \text{ sec.}; x_d1=1.907; x_{q1}=0.278. \)
- G2: \( P_{G2}=200 \text{ MW}; \cos \varphi_{G2}=0.8; U_{G2}=500 \text{ kV}; T_{j2}=5.4 \text{ sec.}; x_d2=1.915; x_{q2}=0.275. \)

Nodes are interconnected by appropriate overhead transmission lines L1-L4.
- L1: \( U_{L1}=500 \text{ kV}; \ell_{L1}=195 \text{ km}; r_0=0.0397 \text{ Ohm/km}; x_0=0.31 \text{ Ohm/km}. \)
- L2: \( U_{L2}=500 \text{ kV}; \ell_{L2}=115 \text{ km}; r_0=0.0362 \text{ Ohm/km}; x_0=0.306 \text{ Ohm/km}. \)
- L3: \( U_{L3}=500 \text{ kV}; \ell_{L3}=180 \text{ km}; r_0=0.0397 \text{ Ohm/km}; x_0=0.31 \text{ Ohm/km}. \)
- L4: \( U_{L4}=500 \text{ kV}; \ell_{L4}=175 \text{ km}; r_0=0.0397 \text{ Ohm/km}; x_0=0.31 \text{ Ohm/km}. \)

Load node parameters:
- \( P_{\text{Load}}=150 \text{ MW}; \cos \varphi_{\text{Load}}=0.88; U_{\text{Load}}=500 \text{ kV}. \)

Generators are equipped with automatic excitation regulators that react to the deflection and the first derivative of the angle, as well as to the voltage deviation. Perform the calculation of the steady state and check the positivity of the first minor \( q_{11} \) of the matrix of the quadratic form \( Q \) for generating nodes.

The results of the calculation of the steady state (modules of the voltage nodes and the angles of the load):
\[
\begin{align*}
\delta_1 &= 505, 2e^{i \theta_1} \text{kV}; \\
\delta_2 &= 496, 2e^{i \theta_2} \text{kV}; \\
\delta_3 &= 509, 8e^{i \theta_3} \text{kV}.
\end{align*}
\]

Then, using the obtained data, solve the equation of the Lyapunov function in quadratic form with respect to the above systems of equations (9).
The calculation is carried out with a heavier mode - a gradual increase in the active load power from $P_{\text{load}}=150\text{MW}$ to $P_{\text{loadmax}}=200\text{MW}$, which led to an increase in the load angles of the generators, respectively, to $\delta_1=126^\circ$ and $\delta_2=131.5^\circ$. As seen from Fig. 2, the characteristic of variation of the minor $q_{11}$ of the first generator $G_1$ depending on the angle $\delta$ first approaches the stability limit.

![Fig. 2 Changes of the first diagonal minor $q_{11}$ of the matrix $Q$ of the Lyapunov function in quadratic form depending on the angle $\delta$: 1-characteristic $G_1$; 2-characteristic $G_2$](image)

**Conclusion**

1. The stability of an electrical system of any complexity with small deviations is characterized by the positiveness of the first minor of the matrix of the quadratic form $q_{11}>0$ of the Lyapunov function in the quadratic form. Moreover, it is shown that if $q_{11}>0$, then all other minors of $Q$ are positive. Therefore, the study of the stability of an EPS of arbitrary complexity can be limited to investigating only the condition $q_{11}>0$. We called this criterion simplified [1, 7], since the positiveness of $q_{11}$ determines the positiveness of the remaining higher minors of the matrix, and therefore only this condition is considered. It is important that the condition $q_{11}>0$ contains theoretically known types of violation of the stability of an electrical system under small disturbances (aperiodic violation, self-excitation, self-swinging), therefore, contains both necessary and sufficient conditions for stability.

2. Analytically and computationally, the $i$-th generator of a complex electrical system is the first to approach the stability limit.

According to the authors, studies of small oscillations of an electrical system based on Lyapunov functions in quadratic form should be developed and carried out in the following directions:

- development of an algorithm and a model for optimal control, assessment and synthesis of the corresponding control laws for EPS with the probabilistic nature of the initial information.

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