Cathode Loading Impact on Voltage Cycling Induced PEMFC Degradation: A Voltage Loss Analysis

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This study focuses on voltage cycling induced degradation of cathodes with different loading (0.4 and 0.1 mgPt/cm²) when applying square wave or triangular wave based accelerated stress tests (ASTs) between 0.6 and 1.0 VRHE. The degradation of the H₂O₂ and H₂/air performance upon extended voltage cycling (up to 30000 cycles) was analyzed in terms of the voltage loss contributions from ORR kinetics, O₂ mass transport resistances and proton conduction resistances in the cathode. The extent of cathode thinning due to carbon support corrosion was determined by post mortem electrode thickness measurements. Square waves were found to cause a more rapid loss of ECSA and mass activity compared to triangular waves, which was shown to be due to the longer hold periods at high potentials rather than to the rate of the potential transient. The observed increase of the O₂ mass transport resistance with voltage cycling was found to mainly depend on the available Pt surface area, while mass transport resistances due to carbon corrosion were found to be insignificant. Finally, it was shown that by lowering the upper potential limit to 0.85 VRHE, low-loaded catalyst layers can sustain 30000 potential cycles without degradation of the H₂/air performance.

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Currently, proton exchange membrane fuel cells (PEMFCs) are on the verge of widespread commercialization as an alternative power source for the automotive market, impressively demonstrated by the launch of PEMFC based vehicles by Toyota,1 Hyundai,2 and Honda3 in recent years. The major hurdle for a significant market penetration of PEMFC vehicles in the near future is the reduction of system cost and the amount of platinum, while meeting long-term durability targets.4,5 Despite remarkable achievements in lowering the Pt loading of the cathode catalyst layer in state-of-the-art membrane electrode assemblies (MEAs) down to 0.1 mgPt/cm²,6–9 recent studies have shown that additional overpotentials occur at low loadings and high current density.8,10 Even though the exact origin of these losses is subject to ongoing discussions in the literature, its occurrence was directly correlated to the available Pt surface area in the cathode, i.e., to the cathode roughness factor (rf, defined as Pt surface area per geometric surface area in units of cm⁴/m²), so that it is often referred to as Pt specific O₂ mass transport resistance (R¹O₂).8,10 Additionally, it is well-known that MEAs degrade under the dynamic load-cycles relevant for automotive applications due to the associated voltage cycles. This made extended voltage cycling tests a frequently applied accelerated stress test (AST) to probe the durability of the cathode catalyst layer. Despite various reports in the literature employing different aging protocols, such as applying square wave (SW) potential perturbations or constant high potential phases, the most commonly applied aging protocol is based on triangular wave (TW) potential perturbations between cell voltages of 0.6 and 1.0 V at a scan rate of typically 50 mV/s.12–21

Thermodynamically, the carbon black support for the noble metal catalyst in the MEA is not stable under fuel cell operating conditions, since the equilibrium potential of the carbon oxidation reaction (COR) is only 0.2 V_RHE (considering CO₂ as reaction product).22 Fortunately, the sluggish kinetics of the COR prevent the oxidation of carbon at potentials < 0.9 V_RHE, i.e., within the typical fuel cell operating range.17,12,23 and therefore enable its use in the cathode catalyst layer. On the other hand, high voltage periods, occurring for example during system start-up or shut-down, can oxidize the carbon structure in the catalyst layer, eventually leading to a loss of void volume and a thinning of the electrode (often referred to as “cathode thinning”).25–27 Since a sufficiently large void volume is crucial for the transport of O₂ through the cathode catalyst layer, significant corrosion of the carbon support causes an increase of the O₂ mass transport resistance and a decrease of PEMFC performance. In addition, carbon corrosion was found to cause a detachment of Pt from the carbon support, leaving the Pt nanoparticles electrically disconnected from the electrode, which decreases the available electrochemical surface area (ECSA) of Pt. Another source of ECSA loss during ASTs is repetitive changes of the cathode potential, leading to oxidation and reduction of the Pt surface.13,14,20 Upon its reduction/oxidation, Pt surface atoms can diffuse from the nanoparticle into the electrolyte as ionic species, eventually redepositing on existing Pt nanoparticles via a process referred to as Ostwald ripening, resulting in an increase of the average Pt particle size and an associated ECSA loss.21,22 Furthermore, as long as the electrode is held at high potentials, Pt ions can diffuse in the electrolyte phase towards the membrane where they can react with cross-over hydrogen to form electrically insulated Pt deposits within the membrane phase.30 Since these Pt deposits can no longer contribute to the electrochemical processes in the cathode catalyst layer, a significant decrease of the ECSA is observed.

To gain a better understanding of the aging processes occurring during voltage cycling ASTs, high-loaded (0.4 mgPt/cm²) and low-loaded (0.1 mgPt/cm²) MEAs were exposed to different potential cycling profiles: square waves, triangular waves, and triangular waves with an upper and lower potential hold (TW–H). The associated losses of ECSA and ORR mass activity were quantified by cyclic voltammetry and by differential flow H₂/O₂ performance measurements, respectively. Furthermore, the performance degradation of differential flow H₂/air polarization curves was analyzed by quantifying the proton conduction resistance via electrochemical impedance spectroscopy (EIS) and by determining the oxygen transport resistance via limiting current diagnostics. The comparison of MEA performance vs ECSA over the course of extended voltage cycling provides insights into the voltage cycling induced increase of the O₂ transport resistance. In addition, by deliberately corroding the carbon support through a high-potential hold of the cathode, where substantial carbon corrosion occurs (1.2 V_cell for 8 h), resulting in pronounced cathode thinning at a
reasonable ECSA retention, we were able to distinguish between the impact of carbon corrosion and ECSA loss on the limiting current measurements. Comparing the various voltage cycling protocols, it could be shown that the higher degradation observed for square wave compared to triangular wave modulation were due to the extended time at high potentials in the former rather than to the voltage ramp rate between upper and lower potential. Our analysis also showed that the oxygen transport resistances developing over the course of voltage cycling ASTs predominantly depends on the specific current density (in units of A/cm²), independent of the initial Pt loading of the MEAs.

Experimental

Comprehensive experimental details on MEA preparation and fuel cell testing are given in our previous publication, but the most pertinent information is given below. All potentials referred to in cyclic voltammograms and in voltage cycling experiments are referenced to the reversible hydrogen electrode (RHE) potential at a nominal H₂ pressure of 100 kPaabs (VRHE). This was obtained by correcting for the Nernstian shift of the hydrogen anode if the H₂ partial pressure was different from 100 kPaabs. For the voltage cycling tests conducted under H₂/N₂ (anode/cathode) at 100 kPaabs, 80°C, 100% RH, the different applied potential limits of 0.60, 0.85, and 1.00 VRHE correspond to cell voltage limits of 0.59, 0.84, and 0.99 Vcell. If not stated otherwise, area-normalized currents (e.g., cm²) refer to the geometric area of the electrode. Furthermore, note that all pressures were measured and controlled at the inlet of the fuel cell and are stated as absolute pressures.

Membrane electrode assembly preparation.—All fuel cell tests were conducted with 5 cm² MEAs, fabricated by the decal transfer method. Catalyst inks were prepared by mixing the catalyst with a low equivalent weight ionomer in a water-solvent dispersion (Asahi Kasei, 700 EW). To realize the same electrode thickness (identical carbon loading of 0.4 mgg/cm²), two catalysts with different Pt/C ratio were used. For low-loaded catalyst layers, a 20%Pt, Pt/Vulcan catalyst (TEC10V20E, Tanaka Kikinzoku Kogyo K.K., TKK) and for high-loaded catalyst layers, a 50%Pt, Pt/Vulcan (TEC10V50E, TKK) was used. The ink components were added into a 15 mL capped bottle (HDPE), already containing 16.5 g of 5 mm ZrO₂ beads as grinding medium in the following sequence: catalyst, water, 1–propanol, and finally the ionomer dispersion. The weight concentration of the ink was 10%, while the solid content was 0.03 g/mLink to obtain a suitable viscosity for the coating process. The ionomer to carbon weight ratio (I/C) was 0.65. The inks were mixed by placing the bottles onto a roller-mill (60 rpm) for 18 h at room temperature. Thereafter, the inks were coated on virgin PTFE using the mayer rod technique with the appropriate bar on a coating machine. The same type of anode was used for all experiments, namely 0.1 mgg/cm² consisting of 20%Pt, Pt/Vulcan (TEC10V20E, TKK) with an I/C ratio of 0.65. The MEAs were assembled by hot pressing a 15 μm membrane (Asahi Kasei) placed between the anode and cathode decals at 155°C for 3 min with an applied force of 0.11 kN/cm². The cathode loadings were 0.1 mgg/cm² and 0.4 mgg/cm², with a maximum deviation of ±5%, as determined by weighing the decals before and after the catalyst layer transfer. For each aging procedure, two independent fuel cell measurements were conducted and subsequently averaged with error bars corresponding to the mean absolute deviation.

Fuel cell testing equipment.—Electrochemical measurements were performed on an in-house manufactured, single-cell hardware, using commercial graphite flow fields (0.5 mm lands/channels; manufactured by Poco Graphite according to our design). Gas diffusion layers (GDLs) were supplied by Freudenberg (H14C7) and the GDL compression was adjusted to 20 ± 1% by PTFE coated, incompressible fiberglass gaskets, assembled at a torque of 12 Nm (for details see Simon et al.). Fuel cell tests were performed on automated test stations (G60, Greenlight Innovation) equipped with a potentio-

![Figure 1. Overview of the experimental procedure for MEA aging. Complete MEA characterization (differential flow H₂/Ω₂ and H₂/air curves, cyclic voltammetry, CO stripping, limiting current measurements, and EIS) was performed at beginning-of-test (BOT) and end-of-test (EOT). The color of the boxes indicates the type of MEA characterization after specific MEA aging cycles: after every aging cycle (black); only at BOT and EOT (orange); after 10000 and 20000 cycles (blue), after 10000 cycles (green).](image-url)
3.5 mV between 500 kHz and 0.2 Hz (20 points per decade), according to a procedure reported in the literature,35,36 based on a transmission line model. In order to ensure reproducibility, the measurement was repeated three times at each condition: 100, 70, 50, and 30% RH, while maintaining constant gas partial pressures (i.e., at cell pressures of 270, 255, 246, and 236 kPaabs, respectively) under differential flow conditions (H2/N2 at 1000/1000 nccm) at 80 °C. The proton resistivity, Rpro (in units of Ω cm), was calculated by dividing Rpro in the cathode thickness.

Differential flow polarization curves were recorded in current-control mode at 80 °C, 170 kPaabs inlet-temperature pressure, 100% RH for both reactants, and constant flows of 2000 nccm of H2 on the anode and 5000 nccm of O2, air, or 10% O2 (balanced with N2) on the cathode. Anode and cathode were operated in counter flow mode. At these conditions, the inlet to outlet pressure drop in anode and cathode was about 2 and 22 kPaabs, respectively. Prior to recording the ASTs were controlled by a potentiostat (Reference3000, Gamry). To avoid reductive currents during potential cycling, all residual O2 was removed from the humidifier and gas line system by a N2 purging procedure, lasting about 45 min. The scan rate for all TW procedures was 50 mV/s, while the potential step in the SW procedure was completed in less than 0.1 seconds (corresponding to an average rate of > 5 V/s). The hold times of 8 s at the vertex potentials in the SW procedure was chosen such that it is identical to the overall time required for one TW cycle (compare Figures 2a and 2b).

Electrode thickness determination by SEM.—The thicknesses of pristine (BOT), as well as aged MEAs (EOT) were determined by cross-section images, obtained by scanning electron microscopy (SEM). To get a more accurate comparison between BOT/EOT and to exclude any effects of cell assembly on the catalyst layer thickness, the pristine MEAs were aged in a cell and operated at 80 °C, 100% RH and 500 nccm N2/N2 flows at ambient pressure for 24 h before determining their initial thickness. The MEAs were cut in quarters, using one quarter close to the gas inlet and one close to the outlet of the cell. The samples were fixed in a Teflon holder and embedded in epoxy resin (EpoThin 2 resin and hardener, Buehler Ltd.) under vacuum pressure to exclude any gas inclusions. After hardening overnight at 40 °C, the resin block was removed from the holder, ground on SiC paper (CarbiMet S, P280, Buehler Ltd.) under constant water flow and thoroughly cleaned by ultrasonification in water. SEM (JCM-6000PLUS NeoScope, Jeol) images were taken in backscattering mode at 15 kV accelerating voltage and a magnification of 2000x at ten arbitrarily chosen locations for each quarter of an MEA, generating 20 images per MEA. The thickness of the electrode was determined by integration of the electrode area with ImageJ (version 1.51j8) and dividing by the length of the image as shown in Figure 3. This way, the entire image is evaluated, which provides a more accurate measure of the catalyst layer thickness compared to a single point measurement.

The thickness of pristine electrodes was used to calculate the packing density of the Vulcan carbon based electrodes. The obtained packing density of 22 ± 4 μm/(mg C/cm2) is in reasonable agreement with values in the literature.37 The BOT thickness for all tested MEAs was approximately 10 μm.
calculated from the transferred weight of carbon in the hot pressing process and the packing density.

Results and Discussion

In this study, we present the aging behavior of MEAs exposed to different voltage cycling protocols, with the respective voltage transients shown in Figure 2. The most commonly used AST is a potential scan between 0.6 and 1.0 V_{RHE} at a scan rate of 50 mV/s and 80 °C in inert gas atmosphere, termed TW in this manuscript (Figure 2a). In addition, many researchers reported the aging behavior of MEAs exposed to square wave (SW) voltammetry (Figure 2b), since the combination of fast potential transients and hold periods has proven to accelerate catalyst layer degradation. However, it is not well-understood yet, whether the holding time at the respective potential or the fast transient from one to the other potential is the dominating factor for the observed electrode degradation. To deconvolute between TW and SW modulations, triangular scans were combined with a potential hold period (TW–H) in order to eliminate the fast potential step while maintaining the effect of the hold time (Figure 2c). Furthermore, electrode aging by a TW modulation between 0.6 and 1.0 V_{RHE} will be compared to a TW modulation with the same lower voltage limit but with a lowered upper potential limit of 0.85 V_{RHE} (TW-LUPL), where carbon support corrosion can safely be assumed to be negligible and where Pt dissolution rates are expected to be reduced (Figure 2d). Finally, our study is complemented by a potential hold experiment at 1.2 V_{cell} for 8 h (under N_2 at 95 °C, 100% RH, and 100 kPaabs) to degrade the carbon support structure via COR, while maintaining a reasonably high electrode roughness factor. All of the aging procedures were conducted on high-loaded (0.4 mgPt/cm^2) and low-loaded cathodes (0.1 mgPt/cm^2) to identify similarities and/or differences in the aging behavior with respect to the ECSA and mass activity originating from different catalyst layer loadings. Moreover, changes of the cathode O_2 mass transport resistance upon aging, especially with respect to the local O_2 mass transport resistance which is commonly observed to be a function of the electrode r_f was of major interest in this study. The utilization of high- and low-loaded catalyst layers enables a comparison over a wide range of r_f, i.e., to compare low Pt specific current densities (A/cm^2) for high-loaded MEAs (high r_f) to high specific current densities after aging of low-loaded cathodes (low r_f). However, to establish a valid comparison of the total mass transport resistance, the same catalyst layer thickness was maintained (≈10 μm) for all MEAs at BOT by utilizing a 50% Pt/C catalyst for high-loaded and a 20% Pt/C catalyst for low-loaded electrodes.

ECSA evaluation: HUPD vs CO stripping.—To evaluate the degradation of the Pt surface area over the course of MEA aging, CVs were recorded after each aging stage, exemplarily shown in Figure 4a for a single low-loaded MEA (0.1 mgPt/cm^2) exposed to SW aging between 0.6 and 1.0 V_{RHE} (see Figure 2b). To determine the ECSA by integration of the HUPD charge, additional CVs with a limited upper potential (0.6 V_{RHE}) were recorded (not shown) to avoid a contribution of charge originating from the reduction of Pt dissolution/redistribution mechanisms, the HUPD features of the CV degrade significantly, i.e., a decrease of the cathode r_f is observed over the course of extended voltage cycling. In general, the evaluation of the ECSA via integration of the pseudo-capacitive HUPD current is influenced by a superimposed faradic current at low potentials, originating from the hydrogen evolution reaction (HER). This superposition is especially dominant at low r_f, as the rate of the HER above 0 V_{RHE} in cyclic voltammetry is mostly thermodynamically controlled (i.e., it does not depend significantly on the active Pt surface area), whereas the pseudo-capacitive HUPD charge scales directly with the available Pt surface area. Hence, the strong decrease of the r_f in the course of the voltage cycling ASTM renders the determination of the ECSA via the HUPD method increasingly erroneous, requiring the implementation of CO stripping to reliably determine the remaining Pt surface area for very low r_f values, i.e., towards the end of the aging protocol.

Representative CO stripping measurements are shown in Figure 4b after subtraction of the second anodic scan. The peak for CO oxidation occurs at ≈0.81 V_{RHE}, which is consistent with reports by other researchers for CO oxidation on Pt nanoparticles, mainly dominated by CO oxidation on Pt(111) facets. Furthermore, a second peak at ≈0.74 V_{RHE} is observed, representing Pt(110)-like planes. Since those facets are less prone to dissolution upon surface oxidation/reduction, this peak becomes more dominant upon potential cycling. Since ECSA determination via HUPD is reasonably accurate for high r_f values, the charge ratio between HUPD and CO stripping remains constant (ECSAHUPD/ECSACO ≈ 0.8) throughout the first phase of aging (up to ≈3000 cycles). As the H_{2} formation is not quantitative in the applied potential range (≥0.07 V_{RHE}), the ECSA determined by the HUPD method is slightly underestimated compared to the ECSA determined from CO stripping, hence their ratio is below unity. In the course of potential cycling, the ratio decreases significantly (ECSAHUPD/ECSACO < 0.1 after 10000 cycles), as the electrode r_f approaches very low values where the HUPD currents are small compared to the capacitive current contributions from the carbon support. As CO stripping was only implemented for electrodes with a considerably low r_f, and HUPD was the preferred method for ECSA determination (due to its experimental simplicity), comparability of these two techniques was assured by introducing a scaling factor. The scaling factor was calculated as the ratio of ECSA determined by HUPD and CO stripping at BOT and was used to correct the values determined by CO stripping to an “HUPD-derived” ECSA at low r_f (denoted by open symbols in figures showing ECSA data in this manuscript).

ECSA evolution during voltage cycling ASTs.—A decrease of the ECSA with increasing number of voltage cycles for each cathode loading and aging protocol is given in Figure 5, eventually reaching similarly low ECSA values on the order of 5–10 m^2/g after 10000 cycles. The ECSA evolution was determined by CO stripping at 0.6 V_{RHE} and compared to HUPD at 0.81 V_{RHE} for high-loaded and 0.74 V_{RHE} for low-loaded cathodes, respectively.
30000 cycles (apart from the TW-LUPL and CP protocols). Since the change of potential in the applied ASTs results in a repetitive reduction/oxidation of the Pt surface, Ostwald ripening and a loss of Pt into the ionomer phase are expected to be the major contributors to the overall loss of active surface area. The gradual stabilization of the ECSA after extended voltage cycling, also observed in previous studies, is likely due to the higher thermodynamic stability of larger Pt particles against dissolution. In contrast to the similar trend of degrading ECSA for all tested samples, its evolution with respect to the aging protocol and the cathode loading differs widely. First of all, it is noteworthy that the ECSA at BOT is higher for low-loaded cathodes (68 ± 1 m^2/g) compared to those with a higher loading (55 ± 1 m^2/g). As mentioned earlier, the electrode thickness was maintained constant for all tested MEAs (≈10 μm), which required a catalyst with 20%Pt content to prepare low-loaded catalyst layers, whereas 50%Pt/C was used for high-loaded electrodes. Even though both catalysts were obtained from the same supplier, Pt nanoparticles are commonly more homogeneously dispersed on the carbon support at low Pt to carbon ratio. This leads to a lower degree of Pt agglomeration, hence a lower average particle size and higher ECSA.\(^{50}\) In fact, an evaluation of the particle size distribution (Figure 5c) by counting ≈400 individual particles (in TEM images) revealed a number averaged diameter of 2.8 ± 0.8 nm and 3.2 ± 1.0 nm for 20%Pt and 50%Pt/C, respectively. As reported frequently by other researchers, Ostwald ripening is more pronounced for small particles compared to larger ones due to their intrinsically lower stability versus dissolution upon oxidation/reduction.\(^{39}\) In accordance with this, for any given aging protocol, a faster decay of the ECSA is observed for low-loaded (orange colored in Figure 5) compared to high-loaded cathodes (blue colored in Figure 5). In summary, the observed relatively stronger ECSA fade for lower catalyst layer loadings in this study is not primarily caused by the cathode loading itself, but is most likely due to the different aging behavior of Pt nanoparticles with different size and different degrees of agglomeration.

Comparing different aging protocols, the ECSA loss is more pronounced for MEAs subjected to SW aging (long-dashed lines) compared to TW aging (solid lines), resulting in a dramatically low ECSA already after 10000 square wave cycles between 0.6 and 1.0 V_RHE (1 m^2/g for high-loaded and 5 m^2/g for low-loaded cathodes). The accelerated degradation caused by square wave compared to triangular wave aging has frequently been observed for PEMFCs,\(^{42}\) as well as in liquid electrolyte.\(^{46}\) There are four possible hypotheses to rationalize the faster degradation during SW cycling: i) a more significant carbon corrosion due to an extended hold period at high potential, eventually resulting in Pt particle detachment, accompanied by a thinning of the cathode catalyst layer; ii) the fast potential transient of the square wave modulation (>4 V/s) could induce more severe Pt dissolution forced by a fast reconstruction of the Pt surface from an oxidized to a reduced state; iii) the holding time at high potential induces a growth of the oxide, resulting in a stronger dissolution in the subsequent cathodic scan; or, iv) the hold time at high anodic potentials provides more time for Pt dissolution and diffusion through the ionomer phase. As shown by Cherevko et al., polycrystalline Pt in 0.1 M H_2SO_4 dissolves at constant anodic potential during about 600 s which is in the same order of magnitude as the dissolution during a potential transient.\(^{52}\) The constantly released Pt ions can diffuse towards the membrane, where they are reduced by H_2 permeating through the membrane from the anode, causing a redeposition of electrically disconnected Pt particles within the ionomer phase.\(^{20}\) The following analysis will shed some light onto the origin of the different aging of MEAs exposed to SW or TW potential profiles as well as at the degradation caused by holding at high potential to deliberately corrode the carbon support.

To investigate whether significant carbon support corrosion takes place over the course of the various voltage cycling protocols, which should lead to cathode thinning, MEA cross sections were prepared at EOT and imaged by SEM. Since the initial thickness of the electrodes at BOT could not be determined by SEM without damaging the MEA, the packing density of high-loaded and low-loaded catalyst layers was evaluated by measuring the thickness and catalyst loading of pristine MEAs (amounting to 22 ± 4 μm (mgPt/cm^2), see experimental section), from which the initial thickness of each MEA could be calculated by its measured catalyst weight. The catalyst layers exposed to TW or SW aging between 0.6 and 1.0 V_RHE (Figures 6a/6b/6d/6e) show no catalyst layer collapse compared to a pristine MEA (Figure 6c), eliminating carbon corrosion as a likely contributor to the overall degradation. Interestingly, the catalyst layers depicted in Figures 6a/6b and 6d/6e show a bright area at the cathode/membrane interface, representing an accumulation of Pt in the membrane phase, caused by precipitation of dissolved Pt in the ionomer phase.\(^{20,30}\) However, due to the rather qualitative nature of such SEM images, no conclusion on the quantitative differences between the amount of deposited Pt during TW and SW cycling can be drawn, even though the larger ECSA loss in the SW protocol points towards a higher amount of disconnected Pt. The absence of cathode thinning was also confirmed for all other
Despite the large extent of carbon corrosion, the reasonably large value of $\eta_{TS}$ might occur (small enough to not lead to cathode thinning), we cannot fully exclude the possibility that a small extent of carbon corrosion or the fast voltage transient. It shall be mentioned that the aging per time interval was in fact smaller for TW-H compared to SW, since both methods showed a similar degradation per cycle while one scan in the TW-H procedure required double the time (32 sec) compared to SW (16 sec). However, the cathode potential remains below 0.9 $V_{RHE}$ for a substantial fraction of the time, while we believe that higher potentials are most detrimental for the ECSA.

To probe the degradation at lower potentials, a triangular scan procedure with a lower upper potential limit (TW-LUPL) of 0.85 $V_{RHE}$ (orange star symbols and short-dotted lines in Figure 5) was tested. In this case, carbon corrosion most certainly will not play a role, so that aging from Pt detachment from the carbon support can surely be excluded (the above analysis suggests, carbon corrosion is not a major contributor to ECSA loss even for upper potentials as high as 1.0 $V_{RHE}$). In addition, Pt dissolution/redeposition effects should also decrease substantially, since the degree of Pt surface oxidation is limited, leading to less Pt dissolution, hence to a smaller overall ECSA loss (40% after 30000 aging cycles). Since the ECSA loss during TW-LUPL is comparatively low, we conclude that most of the additional time during TW-H compared to SW is non-damaging, e.g., below 0.85 $V_{RHE}$. We therefore believe that the lower amount of ECSA loss per time (not per cycle) in the TW-H procedure can be mainly attributed to the additional time at these low, non-damaging potentials. In the light of the rather small ECSA loss when limiting the upper potential to 0.85 $V_{RHE}$ (TW-LUPL) compared to the TW aging protocols with an upper potential limit of 1.0 $V_{RHE}$ for which ECSA losses of 90% over 30000 cycles are observed, avoiding high voltages (e.g., OCV) appears highly advisable to maximize PEMFC durability.

**AST induced ORR activity changes.**—Assuming simple Tafel kinetics for the ORR according to Equation 1, lowering the cathode $\eta$ over the course of ASTs should directly cause an increase of the ORR overpotential according to $^{6,54}$

$$\eta = TS \cdot \left[ \lg \left( \frac{i}{i_0} \right)^{1/\eta_{TS}} \right]$$

where $\eta$ denotes the ORR overpotential in mV, $T$ is the Tafel slope in mV/dec, $i$ denotes the current density in A/cm$^2$, $i_0$ is the exchange current density in A/cm$^2$, and $\eta_{TS}$ denotes the roughness factor in cm$^2$/g.

In Figure 7a, differential flow H$_2$/O$_2$ polarization curves at 80 °C and 100% RH are shown for all AST protocols at BOT and after 30000 aging cycles. First of all, a voltage offset is observed in the initial polarization curves of low-loaded (solid orange lines) compared to high-loaded (solid blue lines) cathodes at BOT. Based on simple Tafel
kinetics (Equation 1) and an ORR Tafel slope of 70 mV/dec, four times lower Pt cathode loadings should result in an approximately four times lower roughness factor (as the catalyst ECSA is very similar) and thus to an additional overpotential of ≈42 mV at the same current density, which is in excellent agreement with the observed offset between high- and low-loaded MEAs (e.g., 737 ± 3 mV vs 695 ± 8 mV at 1.5 A/cm² (42 ± 11 mV) for high- and low-loaded electrodes, respectively).

Over the course of 30000 voltage cycles (≈ EOT condition), the polarization curves of all tested MEAs degrade significantly, with a more severe performance decay for low-loaded MEAs (e.g., 214 mV loss at 2 A/cm²) compared to high-loaded MEAs (39 mV loss at 2 A/cm²). When comparing polarization curves from different aging protocols, the previously identified higher ECSA loss for SW compared to TW cycling correlates well with the observed voltage losses have also been observed by other researchers.9,10,55 Since this overpotential occurs at high geometric current densities and low cathode rf values, i.e., at high local oxygen flux to the Pt surface, we hypothesize that it is related to oxygen transport limitations at the catalyst/ionomer interface. Suggestions for the observed voltage losses at high Pt-specific currents by other research groups include a limited dissolution of oxygen into the ionomer phase, an intrinsically high transport resistance of the thin ionomer phase, a limited effective proton conduction resistance in the cathode catalyst layer, and thus to an additional overpotential of ≈42 mV at the same current density (near 0.9 VRHE), increases to 79 ± 1 mV/dec for TW aged MEAs and more significantly (≈≈90 mV/dec) for the SW and TW–H aging protocols. Tafel slopes of around 120–160 mV/dec were observed for low-loaded MEAs subjected to 30000 TW and SW cycle protocols, while the Tafel slope remained unchanged over the course of the TW–LUPL voltage cycling AST. An increasing Tafel slope due to loss or restructuring of the ionomer can be excluded as the proton transport resistance which apparently increase with extended voltage cycling. Since the electrode thickness remained unchanged (Figure 6f), even for the most damaging aging protocols (SW and TW–H), a hindrance of the gas phase mass transport caused by a deterioration of the overall catalyst layer structure can be excluded. Although, slight corrosion of the catalyst support and a concomitant slight decrease of the available pore volume cannot be fully observed, cathode thinning was clearly not observed by SEM in any of our experiments.

Additionally, changes of the Tafel slope due to loss or restructuring of the ionomer can be excluded as the proton transport resistance did not change significantly upon voltage cycling (not shown) and major chemical degradation of the ionomer can be excluded due to the absence of oxygen in the largest fraction of the experimental time.56

We would like to emphasize that these measurements were performed with pure O₂, where gas phase O₂ transport resistances can be excluded as possible cause for the increasing apparent Tafel slopes. It is also worth noting that these severe and unassigned mass transport overpotentials occur at similar values of the Pt specific current density (near ≈≈0.05 A/cm²) in all H₂/O₂ polarization curves. At high Pt specific current densities, caused by ultra-low rf values, additional voltage losses have also been observed by other researchers.9,10,55,56 Since this overpotential occurs at high geometric current densities and low cathode rf values, i.e., at high local oxygen flux to the Pt surface, we hypothesize that it is related to oxygen transport limitations at the catalyst/ionomer interface. Suggestions for the observed voltage losses at high Pt-specific currents by other research groups include a limited dissolution of oxygen into the ionomer phase, an intrinsically high transport resistance of the thin ionomer phase, a limited effective ionomer surface, or a decrease of the ORR kinetics at low potential.59 Although we cannot exclude any one of these effects, a Pt/ionomer specific resistance appears most reasonable to explain the observed voltage losses at high Pt-specific currents, as the voltage profile shows a shape resembling a mass transport limitation, rather than a kinetic phenomenon.

A linear regression of the measured data in the current density range from 50 to 500 mA/cm² and extrapolation to 0.9 V_RHE yields the kinetic ORR activity at 0.9 V_RHE (at 120 kPa abs, O₂ partial pressure), commonly used to compare the catalytic activity of platinum. It can either be normalized to the measured ECSA to yield the specific current density for the ORR at 0.9 V_RHE (i^*val, Figure 8b) or to the mass of Pt in the cathode catalyst layer to yield the mass activity for the ORR at 0.9 V_RHE (i^*val, Figure 8a) to enable comparability between all experiments. In general, the here determined ORR
TW–LUPL procedure (orange stars) is significantly lower compared during square wave aging. Finally, the ORR mass activity loss over the potential hold phase is the dominating factor for MEA degradation (see Figure 5), underlining our previous hypothesis that identical to those aged by SW cycles (compare blue triangle and diamond symbols in Figure 8a), the decrease of imass to the stronger lines, diamonds), the decrease of imass as a function of cycle number is more pronounced for SW profiles, which is directly correlated to the other protocols, which again is in accordance with the lower ECSA loss.

It is interesting to note that in contrast to the overall trend of decreasing ORR mass activity, i_{\text{ORR}}^{\text{max}} at 0.9 V reach a maximum after 100 cycles, particularly noticeable for low-loaded MEAs. This could be ascribed to an activation of the catalyst particles due to cleaning of the Pt surface, which might overcompensate the ECSA loss of \( \approx 20\% \) in the initial 100 cycles. If true, this would mean that the initial activation procedure (see experimental section) was not sufficient to fully activate the cathode catalyst layer. Alternatively, there is the possibility that the measurement sequence (described in Figure 1) might influence the observed small variation of the measured ORR activity. While the O

The ORR mass activities were always higher when measured directly after a recovery step (i.e., at 100, 1000, 3000 and 20000 cycles), while reproducibly lower values were obtained in the absence of a directly preceding recovery step (at BOT as well as at 10000 and 30000 cycles). Since the overall degradation rate induced by the TW-LUPL procedure is rather small (\( \approx 10\% \) loss in ECSA and \( i_{\text{ORR}}^{\text{max}} \) after 30000 cycles), this presumably testing procedure related inaccuracy of \( \approx 10\% \) in the quantification of \( i_{\text{ORR}}^{\text{max}} \) might have become apparent. Nevertheless, this minor uncertainty in ORR mass activity values is of no consequence to the conclusions drawn from these data.

Further insights into the aging process can be obtained by examining the specific ORR activity at 0.9 V vs RHE, \( i_{\text{ORR}}^{\text{spec}} \), as shown in Figure 8b. It is well-known that small Pt nanoparticles have a lower specific ORR activity due to their relatively larger amount of edge and corner sites compared to larger nanoparticles.\(^{60}\) Amongst others, Mayrhofer et al. found an increase of the specific activity of Pt particles by a factor \( \approx 2 \) when going from 1 nm sized particles to 5 nm.\(^{60}\) In our case, the determined specific ORR activities of low-loaded (182 \( \pm \) 18 \( \mu \)A/cm\(^2\)) and high-loaded MEAs (177 \( \pm \) 16 \( \mu \)A/cm\(^2\)) at BOT are essentially identical, consistent with the very similar average particle size of the two catalysts (2.8 \( \pm \) 0.8 nm for 20\%Pt/PtC and 3.2 \( \pm \) 1.0 nm for 50\%Pt/PtC). In the course of aging, the specific activity increases for all tested MEAs due to the constantly growing fraction of larger particles. The highest observed ORR activity (\( \approx 1 \) mA/cm\(^2\)) was found for low-loaded catalyst layers with the most significant ECSA loss (Figure 5) due to the harsh SW aging protocol. However, it shall be noted that the determination of the catalyst activity is slightly impaired by an increased Tafel slope for MEAs with a very small \( rf \) value, hence the calculation of \( i_{\text{ORR}}^{\text{spec}} \) becomes increasingly erroneously for strongly aged catalyst layers due to difficulties in determining kinetic parameters accurately.

**AST induced H\_2/air and H\_2/10\%O\_2 performance losses.**—To obtain further information about mass transport related overpotentials induced by the different aging protocols, polarization curves in air and 10\% O

The TOD mass activities at BOT (\( \approx 110–140 \) A/gPt for the 20\%wt Pt/V catalyst and \( \approx 85–105 \) A/gPt for the 50\%wt Pt/V catalyst) compare well with those reported by other researchers for Pt supported on Vulcan type carbon.\(^{60,61}\) In accordance with the severe ECSA loss in most of the here examined aging protocols, the overall surface available for the ORR decreases significantly from BOT to EOT (low-loaded TW, SW and TW–H: \( \approx 90\% \) ECSA loss), there is also a significant loss of \( i_{\text{ORR}}^{\text{max}} \) over the course of aging for all loadings and aging procedures (see Figure 8a). Comparing TW (solid lines, circles) and SW (long-dashed lines, diamonds), the decrease of \( i_{\text{ORR}}^{\text{max}} \) with increasing \( \rho \) from 10 mV vs 55 mV at 1.5 A/cm\(^2\) for high- and low-loaded electrodes, respectively) and 10\% O\_2 (to simulate conditions at the cathode outlet of a PEMFC stack operated at an air stoichiometry of \( \approx 1.75 \)) were measured at several aging stages (Figure 9). The solid lines are the average BOT performance of MEAs with 0.4 mgPt/cm\(^2\) (blue) and 0.1 mgPt/cm\(^2\) (orange) cathode loadings at the here investigated operating conditions of 170 kPa\sub{atm}, 80 C, and 100\% RH (error bars represent the standard deviation across all measured MEAs, indicating the excellent reproducibility of the performance at BOT). Similarly to polarization curves in O\_2 (Figure 7a), a clear offset between high-loaded and low-loaded cathodes, resulting from a different cathode \( rf \) is observed for BOT polarization curves in air (e.g., 689 \( \pm \) 3 mV vs 634 \( \pm \) 10 mV at 1.5 A/cm\(^2\) for high- and low-loaded electrodes, respectively) and 10\% O\_2 (e.g., 620 \( \pm \) 10 mV vs 496 \( \pm \) 9 mV at 1.5 A/cm\(^2\) for high-
High-loaded MEAs aged by the SW and TW-H protocol (open blue diamonds and triangles in Figure 9) show a similar performance at EOT, which was significantly lower compared to MEAs aged by the TW protocol (open blue circles). This is in accordance with the above discussed O₂ polarization curves (Figure 7a) and ECSA losses (Figure 5a). Even though the overall performance losses are rather dramatic for the low-loaded MEAs, they also show the same increased aging when subjected to SW rather than TW voltage cycling (open orange diamonds and circles). However, as one might have hoped based on the ECSA (Figure 5) and ORR mass activity (Figure 8a) data, even with low-loaded cathodes no voltage cycling induced performance losses occur over 30000 cycles when using the TW-LUPL protocol with a lowered upper voltage limit of 0.85 V_{RHE}, neither in H₂/air nor in H₂/10%O₂ (orange stars in Figure 9).

An interesting insight into the \( r_f \) dependent \( R^{\text{O}_2} \) induced losses can be gained by comparing the H₂/air EOT performance of the high-loaded cathode subjected to TW cycling (open blue circles) and the low-loaded cathode subjected to TW-LUPL cycling (orange stars). At EOT, the high-loaded cathode has a slightly higher roughness factor (56 ± 2 cm²/cm²) compared to the low-loaded cathode (42 ± 2 cm²/cm²), so that one would expect a similar or an even lower \( R^{\text{O}_2} \) value for the former. At the same time, the specific ORR activity at 0.9 V_{RHE} for the high-loaded cathode (350 ± 29 μA/cm²) is substantially larger than that for the low-loaded cathode (224 ± 7 μA/cm²). Since the product of \( \eta_{\text{ORR}} \cdot r_f \), equating to 19.6 mA/cm² for the high-loaded and 9.2 mA/cm² for the low-loaded cathode, is proportional to \( \eta_{\text{ORR}} \cdot (r_f) \) under the same operating conditions, the \( \eta_{\text{ORR}} \) difference predicted by Equation 1 is ≈23 mV, suggesting that the kinetically predicted EOT performance of the high-loaded cathode subjected to TW cycles should be ≈23 mV higher than that of the low-loaded cathode subjected to TW-LUPL cycles. This is consistent with the higher H₂/O₂ performance of the former (see Figure 7a) and also its superior low current density performance in H₂/air or H₂/10%O₂ (see Figure 9). However, despite its higher EOT roughness factor, the TW-aged high-loaded cathode exhibits a larger (unassigned) mass transport loss, evident by its lower H₂/air or H₂/10%O₂ performance at current densities above ≈1 A/cm² (see Figure 9). Since no significant differences in the effective proton conduction resistance (measured by EIS) were observed between these two electrodes, an apparently larger oxygen transport resistance, which is not simply related to the cathode \( r_f \) value and produced by voltage cycling, must exist. Its origin is unclear at the moment, but it may be related to the morphological changes and redistribution of Pt within the cathode electrode caused by voltage cycling.

\( O_2 \) transport resistance evaluation.—In order to disentangle mass transport resistances originating from gas transport in the flow fields and the GDL from those in the catalyst layer, and to identify resistances dependent on the cathode \( r_f \) value, limiting current measurements were executed according to the method developed by Baker et al. In this approach, the concentration of O₂ (dry gas fraction in N₂) in the feed gas is varied from 4% to 24%, while a low potential between 0.05 and 0.30 V is applied to obtain the limiting current, where the O₂ concentration at the Pt surface decreases to zero and the current-voltage profiles at different O₂ concentrations resemble a vertical line (see Figure 10a). For each concentration, the O₂ mass transport resistance is calculated and plotted versus the respective limiting current density, as illustrated in Figure 10b for the low-loaded cathode MEAs at BOT and after 30000 SW cycles (for details on the analysis see Simon et al.).

For all cathodes at BOT, the total oxygen mass transport resistance, \( R^{\text{O}_2}_{\text{total}} \), is essentially independent of the limiting current density at all pressures (see triangles in Figure 10b and its inset), as is expected under these conditions for the Freudenberg GDL. Furthermore, the determination of the limiting current density at various pressures enables the separation of \( R^{\text{O}_2}_{\text{total}} \) into a pressure dependent \( R^{\text{Pt}}_{\text{O}_2} \) and a pressure independent \( R^{\text{PD}}_{\text{O}_2} \) oxygen mass transport resistance via a linear regression of \( R^{\text{O}_2}_{\text{total}} \) versus the absolute pressure (Figure 10c).

![Figure 9](image-url) Differential flow a) H₂/air and b) H₂/10%O₂ (2000/5000 nccm) performance (170 kPa,gas, 80 C, and 100% RH) for MEAs with 0.4 mgPt/cm² (blue lines/symbols) and 0.1 mgPt/cm² (orange lines/symbols) cathodes at BOT (solid lines and filled symbols) and at EOT, i.e., after 30000 voltage cycles or an 8 hour hold at 1.2 Vcell (dotted lines or hollow symbols). BOT polarization curves of all tested MEAs were averaged and error bars represent the standard deviation of all measurement points. At EOT, error bars represent the mean absolute deviation of two measurements.

low-loaded electrodes, respectively). In contrast to pure O₂, where the offset perfectly matched the increase of \( \eta_{\text{ORR}} \) predicted by Equation 1, the voltage loss for low-loaded MEAs in air (e.g., 55 ± 13 mV at 1.5 A/cm²) and 10% O₂ (e.g., 124 ± 19 mV at 1.5 A/cm²) was much higher than the kinetically expected penalty. The additional voltage loss for low-loaded cathodes in air, being even more pronounced in 10% O₂, is an initial hint on a mass transport overpotential related to the \( r_f \) value of the cathode. Several other researchers reported the occurrence of such a Pt surface area specific O₂ mass transport resistance (termed \( R^{\text{O}_2}_{\text{Pt}} \) in this manuscript) for pristine cathodes with different loadings and showed its substantial increase with decreasing cathode \( r_f \) values by limiting current measurements.\(^{8-10,41}\)
Nevertheless, a severe decrease of the limiting current is observed for low-loaded cathodes after extended TW or SW cycling, exemplarily shown in Figure 10a for the SW protocol (BOT ≡ triangles; EOT ≡ circles). At EOT, the limiting current density (circles in Figure 10a) is significantly lower, and the resulting $R_{O2}^{total}$ values are more than an order of magnitude larger than at BOT and change severely with limiting current densities, especially at higher pressures (circles in Figure 10b). We attribute this unsteady transport resistance to a local catalyst layer/GDL flooding (at an inlet RH of 100%) due to the low limiting current, at which no significant through-plane temperature gradient between MEA and flow field can be established, thereby hampering product water removal through the gas phase. This hypothesis is underpinned by the fact that the transport resistance is strongly decreasing with increasing current density, i.e., at a higher temperature gradient. The effect of flooding is even more pronounced at higher applied pressures, analogous to what is generally observed.

As the evaluation of $R_{O2}^{total}$ at higher pressure is not reliable in cases like the one shown in Figure 10b (circles), a deconvolution of $R_{O2}^{total}$ into $R_{O2}^{PD}$ and $R_{O2}^{PI}$ is not trustworthy, hence only the total transport resistance at 170 kPaabs is considered in these instances. The origin of this effect is yet unclear, but as will be shown in the following, it is not related to carbon support corrosion.

The total oxygen mass transport resistance at different stages of the aging protocol is shown for high-loaded (blue) and low-loaded (orange) cathodes in Figure 11, separated into pressure dependent, $R_{O2}^{PD}$ (solid bars), and pressure independent, $R_{O2}^{PI}$ (hollow bars) contributions, according to the above described analysis. All high-loaded (average of 0.73 ± 0.09 s/cm) as well as all low-loaded cathodes (average of 0.82 ± 0.10 s/cm) show a similar $R_{O2}^{total}$ at BOT, demonstrating a reproducible MEA fabrication. In general, low-loaded cathodes show a slightly higher transport resistance than high-loaded electrodes at BOT, as already indicated by the observed voltage losses in H2/air (Figure 9a) and H2/10% O2 polarization curves (Figure 9b). The pressure dependent transport resistance associated with the GDL is, as expected, essentially identical for high- ($R_{O2}^{PD} = 0.44 ± 0.05$ s/cm) and low-loaded ($R_{O2}^{PD} = 0.45 ± 0.05$ s/cm) cathodes at BOT. Therefore, the differences in $R_{O2}^{total}$ originate from a higher pressure independent transport resistance, due to their lower cathode rf values ($R_{O2}^{PI}$ high-loaded: 0.29 ± 0.04 s/cm; $R_{O2}^{PI}$ low-loaded: 0.37 ± 0.05 s/cm).

Over the course of catalyst layer aging, the pressure dependent transport resistance remains reasonably constant, considering that the observed fluctuations are within the experimental accuracy. At the same time, $R_{O2}^{PI}$ increases significantly. In general, $R_{O2}^{PI}$ is composed of various contributions, namely Knudsen diffusion in small pores, diffusion in the ionomer and in water, as well as of the so-called Pt surface specific resistance ($R_{rf}^{PI}$). Since the ionomer content and the overall electrode thickness was constant in all experiments, we ascribe the higher $R_{O2}^{PI}$ of low-loaded electrodes to a rf specific mass transport resistance. An increase of $R_{O2}^{PI}$ was also observed by other researchers who conducted limiting current measurements for MEAs with different loadings. In accordance with the ECSA trends presented in Figure 5, $R_{O2}^{PI}$ increases with decreasing ECSA values in all aging protocols. Consistently, the highest transport resistances

Figure 10. a) Limiting current measurement in the potential range between 0.30 and 0.05 V at O2 concentrations of 4%, 12%, and 21%, comparing a 0.1 mgPt/cm2 cathode at BOT (triangles) and after 30000 SW cycles (circles). Measurements were performed at 80˚C, 100% RH, and 170 kPaabs (using 2000 nccm H2 and 5000 nccm O2 in N2). b) Total O2 mass transport resistance vs limiting current density at 170, 270, 350 and 500 kPaabs (O2 concentration between 4% and 24%); the inset is a zoom into the low limiting current densities. c) $R_{O2}^{total}$ of the BOT cathode averaged over all O2 concentrations at the respective pressure; the shown linear regression line is used to deconvolute $R_{O2}^{total}$ into $R_{O2}^{PD}$ and $R_{O2}^{PI}$.

Figure 11. Total oxygen mass transport resistance ($R_{O2}^{total}$) determined at 80˚C, 100% RH, and 170 kPaabs (differential flow of hydrogen and 4% - 21% O2 (balance N2) for high-loaded (blue) and low-loaded (orange) cathodes tested by different ASTs (for each, left bars are at BOT, middle bars at 10000 cycles, and right bars 30000 cycles). $R_{O2}^{total}$ was also separated into pressure dependent ($R_{O2}^{PD}$, solid bar) and pressure independent ($R_{O2}^{PI}$, open bar) oxygen mass transport resistance contributions. For certain experiments, the deconvolution of $R_{O2}^{total}$ was not reliable, thus only the total transport resistance at 170 kPaabs is plotted as a dotted bar. Error bars correspond to the mean absolute deviation of two repeat experiments with different MEAs.
were obtained for procedures which showed the most severe ECSA losses, namely SW and TW-H for high-loaded cathodes as well as TW and SW for low-loaded cathodes. Comparing the transport resistance evolution for high-loaded catalyst layers, TW aging leads to less significant increases of \( R_{O2}^{Pi} \) compared to SW and TW-H aging, which in turn show a similar transport resistance at each aging interval. Taking the similar ECSA decrease and the same performance in polarization curves of TW-H and SW into account, it is evident that the potential transient during aging is of minor importance, while the hold time at high potentials is the main contributor to electrode degradation. As described earlier, for degraded cathodes with ultra-low \( rf \) values (low-loaded cathodes subjected to 30000 TW cycles or \( 3 \times 10000 \) SW cycles), the deconvolution of pressure dependent and pressure independent resistance is not reliable, hence only the total transport resistance at 170 kPaabs, is shown in Figure 11.

To disentangle the effect of low \( rf \) values on the observed increase of the total transport resistance from possible transport limitations originating from carbon corrosion, cathodes were exposed to a constant potential hold (CP) of 1.2 \( V_{cell} \) for 8 h at 95 °C and 100% RH. This procedure results in a thinning of the catalyst layer by approximately 40%, while maintaining a reasonably high \( rf \) value (from 222 ± 10 at BOT to 90 ± 2 \( cm^2/Pt \) at EOT). Despite the fact that the catalyst layer structure collapses in this experiment, the total transport resistance increases only moderately from 0.70 ± 0.01 to 0.86 ± 0.01 s/cm (see Figure 11). On the one hand, the slight increase of the pressure independent resistance can be attributed to the loss of ECSA, comparable to the other testing procedures which lead to lower \( rf \) values. On the other hand, since \( R_{O2}^{Pi} \) is mostly associated with the transport resistances in the flow-field, the GDL, and in the void volume within the catalyst layer, it would be reasonable to expect that it increases upon cathode thinning. Nevertheless, since the pressure dependent part of the transport resistance increases solely by 0.07 ± 0.04 s/cm (Figure 11), this hypothesis cannot be proven within the limited accuracy of this method. Considering that no thinning was observed for electrodes exposed to voltage cycling, the changes of \( R_{O2}^{Total} \), and especially \( R_{O2}^{Pi} \), in the course of aging are clearly not attributed to carbon corrosion. In other words, the much larger increase of \( R_{O2}^{Total} \) and \( R_{O2}^{Pi} \) over the course of all voltage cycling protocols (with the exception of the TW–LUPL protocol) where no cathode thinning was observed, must be related to phenomena other than carbon corrosion. Consistent with the essentially identical BOT and EOT performance of low-loaded cathodes subjected to TW-LUPL cycling, there is also essentially no increase of the overall oxygen mass transport resistance.

As a comparison to the data in literature, the pressure independent part of the \( O_2 \) mass transport resistance was plotted as a function of the \( rf \) in Figure 12. The clear trend of strongly increasing \( R_{O2}^{Pi} \) with decreasing platinum surface area (i.e., \( rf < 50 \ cm^2/Pt \)) presented here, was also observed by Greszler et al. and underlines the previously stated hypothesis that the high \( O_2 \) mass transport resistance is a strong function of the cathode \( rf \). The fact that aged, high-loaded electrodes show a similar \( R_{O2}^{Pi} \) (e.g., after 30000 TW cycles) compared to low-loaded cathodes at BOT further clarifies that the \( rf \) is indeed a major influencing factor on the \( O_2 \) mass transport. However, a disentanglement of \( R_{O2}^{Pi} \) into a fraction caused by the AST and another contribution, originating solely from the \( rf \) of the cathode, cannot be drawn unambiguously with respect to the accuracy of the measurement.

Voltage loss analysis of aged cathodes.—Using the total oxygen mass transport resistance at BOT and EOT (Figure 11), the corresponding air polarization curves were corrected for all voltage loss contributions that are currently accessible by in situ measurement techniques, according to a preceding publication and are shown in Figure 13. In short, corrections of the potential include the HFR, the effective cathode proton conduction resistance, and the total oxygen mass transport resistance, while the current density was corrected for \( H_2 \) crossover, electrical short current and was normalized to the cathode’s ECSA.

For BOT, all \( H_2/air \) polarization curves coincide with a Tafel slope of \( \approx 73 \) mV/dec for high-loaded and \( \approx 80 \) mV/dec for low-loaded catalyst layers. We attribute the higher Tafel slope in air compared to pure \( O_2 \) to unaccounted mass transport limitations, which increase as the \( O_2 \) concentration decreases. Upon aging, a strong deviation from the theoretical 70 mV/dec slope is observed at high current density, similar to the limitations observed in pure \( O_2 \) (deviations from a straight \( \tilde{E}_{corrected} \) Vs log(\( i \)) line at EOT occur at \( \approx 30 \) mA/cm\(^2\) for all aged cathodes, except for the low-loaded cathode subjected to SW cycling). However, the deviation from a straight \( \tilde{E}_{corrected} \) Vs log(\( i \)) line at EOT occurs earlier in air, viz., at \( \approx 10 \) mA/cm\(^2\) (again, with the exception of the low-loaded cathode subjected to SW cycling, which deviates much earlier), which one would expect in case of significant oxygen transport resistances. The here observed deviation was also seen by other research groups at specific current densities ranging from 5–50 mA/cm\(^2\) in fully corrected air polarization curves with aged MEAs or at very low cathode \( rf \) values, commonly attributed to a local oxygen diffusion resistance. Although a correction for the mass transport was applied to these measurements, it seems that the deviation at high Pt specific current density is not accounted for by the transport resistance determined by the limiting current method. Indeed, the conditions of the MEA during the limiting current measurement (local MEA temperature and \( RH \) due to a high heat production rate at \(<0.3 \) V) are not identical to the operating conditions during the polarization curve (lower heat production rate at voltages of \( \approx 0.6-0.5 \) V) where the corrections are applied, which certainly introduces an error. Furthermore, due to changes in temperature or local \( RH \), the utilization of the Pt surface area may change, increasing the oxygen flux to the remaining Pt surface or altering the ORR kinetics. While we cannot prove one or the other hypothesis, we consider the limitations of the oxygen mass transport determination method to be a probable cause of at least part of the observed deviations at high current density. However, whatever the origin of these unaccounted transport resistances may be, they clearly scale with the specific current density.

After full correction of the recorded \( H_2/air \) polarization curves for all known voltage loss contributions (namely the HFR, the effective cathode proton conduction resistance, and the total oxygen mass transport resistance, including the pressure independent contributions), the corrected polarization curves for high-loaded cathodes at BOT and EOT further clarify that the HFR, the effective cathode proton conduction resistance, and the total oxygen mass transport resistance, including the pressure independent contributions, are the major loss mechanisms at high current density.
transport resistance), the H\textsubscript{2}/air polarization curves can be compared to the theoretical, kinetically limited polarization curve, constructed from the experimentally determined ORR mass activities (see Figure 8a) and assuming simple Tafel kinetics with an intrinsic Tafel slope of 70 mV/decade. The remaining unaccounted voltage losses are shown in Figure 14a for high-loaded and in Figure 14b for low-loaded catalyst layers at BOT (solid lines) and EOT for all aging protocols.

An analysis of Figure 14a reveals an almost quantitative account of all voltage losses for high-loaded MEAs at BOT (blue solid line), with only minor unaccounted losses at 2 A/cm\textsuperscript{2} (15 ± 11 mV). Upon catalyst layer aging over 30000 TW cycles (open blue circles), the unaccounted voltage losses at higher current density (> 1 A/cm\textsuperscript{2}) increase, pointing to the fact that not all voltage losses can be assessed by the current in situ diagnostic measurements. This discrepancy is even more pronounced for the MEAs exposed to SW and TW-H aging, where unaccounted voltage losses on the order of 250 mV were observed at 2 A/cm\textsuperscript{2}. In accordance with the previously described evolution of the ECSA, the ORR mass activity, and the total oxygen mass transport resistance, both aging protocols (SW and TW-H) lead to identical unaccounted losses. In contrast, the corrosion of the carbon support (CP aging) does not lead to a significant increase of the unaccounted losses (e.g., 41 ± 12 mV at 2 A/cm\textsuperscript{2}) and therefore cannot explain the high, unaccounted voltage losses for SW and TW-H cycling.

The unaccounted voltage losses for low-loaded MEAs (Figure 14b) at BOT are slightly higher compared to high-loaded MEAs, and are in reasonably good agreement with our previous voltage-loss analysis with the same catalyst.\textsuperscript{11} After aging by the TW or SW procedure, rather dramatic unaccounted voltage losses are already observed at low geometric current densities and become increasingly larger with current density. Considering the SW and TW-H aged, high-loaded catalyst layers, the significant unaccounted voltage losses of all these measurements suggest that the quantitative evaluation of voltage loss contributions is impaired when using cathodes with low \(r_f\) values, i.e. at a high Pt specific current density. This hypothesis is further validated by the TW-LUPL voltage cycling test applied to the 0.1 mgPt/cm\textsuperscript{2} cathodes: even after 30000 potential cycles, the unaccounted voltage losses were substantially less compared to all other aging protocols, being as little as 24 ± 1 mV at 2 A/cm\textsuperscript{2} at EOT, where the roughness factor is 42 ± 2 cm\textsuperscript{2}/mg\textsuperscript{Pt}. A very intriguing observation is that this roughness factor is only ≈2–fold larger than the EOT roughness factor of the high-loaded cathodes subjected to SW cycling (29 ± 1 cm\textsuperscript{2}/mg\textsuperscript{Pt}), while the unaccounted voltage losses of the latter are >250 mV at 2 A/cm\textsuperscript{2}. Quite clearly, the currently available in situ diagnostics do not capture the critical voltage loss phenomena which must be related to the morphological changes with regards to Pt particle size and location within the electrode.

Conclusions

To better understand the influence of voltage cycling based ASTs on the development of mass transport resistances and on the extent to which they can be captured by currently available in situ diagnostics (ORR kinetics in H\textsubscript{2}/O\textsubscript{2}, cathode proton conduction resistance by impedance, oxygen transport resistance via limiting current measurements), we conducted a comprehensive voltage loss analysis on MEAs with high- and low-loaded cathodes (0.4 and 0.1 mgPt/cm\textsuperscript{2}) over the course of ASTs based on different voltage cycling protocols. These include triangular wave (TW) and square wave (SW) cycles between 0.6 and 1.0 V\textsubscript{HRE}, TW cycles between 0.6 and a lowered upper potential limit of 0.85 V\textsubscript{HRE} (TW–LUPL), as well as a combination of a triangular sweeps between 0.6 and 1.0 V\textsubscript{HRE}, with potential hold periods at the upper and lower potential (TW–H), all conducted in H\textsubscript{2}/N\textsubscript{2} (anode/cathode) at ambient pressure, 80 °C, and 100% RH. The unaccounted voltage losses at higher current density (0.1 mgPt/cm\textsuperscript{2} cathodes: even after 30000 potential cycles, the unaccounted losses (e.g., 41 ± 1 mV at 2 A/cm\textsuperscript{2}) were further validated by the TW-LUPL voltage cycling test applied to the high-loaded (blue) and low-loaded (orange) MEAs at BOT. BOT polarization curves of all tested MEAs were averaged and error bars symbolize the standard deviation over all MEAs. Error bars for EOT curves represent the mean absolute deviation between two measurements.

![Figure 13. Transport-corrected differential flow H\textsubscript{2}/air polarization curves at 80 °C, 100% RH, and 170 kPa\textsubscript{abs} at BOT (full symbols) and EOT (open symbols) for high-loaded (blue) and low-loaded (orange) MEAs. The potential is corrected for the HFR, effective cathode proton conduction resistance, and the total oxygen mass transport resistance, while the current is corrected for the H\textsubscript{2} crossover and ohmic shorting currents, and is normalized to the cathode ECSA. The two slanted lines indicate the specific current density region for high-loaded (blue) and low-loaded (orange) MEAs at BOT. BOT polarization curves of all tested MEAs were averaged and error bars symbolize the standard deviation over all MEAs. Error bars for EOT curves represent the mean absolute deviation between two measurements.](image1)

![Figure 14. Unaccounted voltage losses between a theoretical, kinetically limited H\textsubscript{2}/air polarization curve calculated from the measured ORR mass activity (assuming simple Tafel kinetics with an intrinsic Tafel slope of 70 mV/dec) and polarization curves recorded in H\textsubscript{2}/air (differential flow, 170 kPa\textsubscript{abs}, 80 °C, and 100% RH), where the potential is corrected for the HFR, effective cathode proton conduction resistance, and the total oxygen mass transport resistance for a) high-loaded and b) low-loaded cathode catalyst layers at BOT (full symbols) and EOT (open symbols).](image2)
A comparison of the ECSA and the O$_2$ mass transport resistance evolution of MEAs aged by SW cycles leads to considerably faster degradation compared to TW cycles. To identify whether the fast potential transients or the potential hold phase is responsible for the faster degradation rates during SW aging, an additional procedure, consisting of a triangular potential scan followed by potential hold phase (TW–H) was tested. A comparison of the evolution of the cathode’s ECSA, mass activity, and mass transport resistance revealed that the potential hold rather than the rate of the potential transient is the main contributor to cathode aging in voltage cycling ASTs, so that SW and TW–H cycling profiles lead to identical aging rates with respect to the number of cycles.

For all ASTs, the magnitude of assignable and unassigned O$_2$ mass transport resistances scaled predominantly with the cathode roughness factor. This analysis showed that the impact of carboxylate groups on the O$_2$ mass transport resistance is minor compared to its increase with decreasing roughness factor. While 30000 TW or SW cycles between 0.6 and 1.0 VRHE dramatically decrease the H$_2$/air performance of low-loaded MEAs, it shows negligible degradation over 30000 cycles when a lowered upper potential limit of 0.85 VRHE is applied (TW–LPL). This provides a guideline to limit cathode degradation under automotive conditions, namely by eliminating high potential periods (e.g., OCV).

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