Stereoscopic measurements of coronal Doppler velocities* 

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ABSTRACT

Context. The Solar Orbiter mission, with an orbit outside the Sun-Earth line and leaving the ecliptic plane, opens up opportunities for the combined analysis of measurements obtained by solar imagers and spectrometers. For the first time different space spectrometers will be located at wide angles to each other, allowing three-dimensional (3D) spectroscopy of the solar atmosphere.

Aims. The aim of this work is to prepare a methodology to facilitate the reconstruction of 3D vector velocities from two stereoscopic line of sight (LOS) Doppler velocity measurements using the Spectral Imaging of the Coronal Environment (SPICE) on board the Solar Orbiter and the near-Earth spectrometers, while widely separated in space.

Methods. We developed the methodology using the libraries designed earlier for the STEREO mission, but applied to spectroscopic data from the Hinode mission and the Solar Dynamics Observatory. We used well-known methods of static and dynamic solar rotation stereoscopy and the methods of extreme ultraviolet (EUV) stereoscopic triangulation for optically thin coronal EUV plasma emissions. We developed new algorithms using analytical geometry in space to determine the 3D velocity in coronal loops.

Results. We demonstrate our approach with the reconstruction of 3D velocity vectors in plasma flows along ‘open’ and ‘closed’ magnetic loops. This technique will be applied to an actual situation of two spacecraft at different separations with spectrometers on board during the Solar Orbiter nominal phase: SPICE versus the Interface Region Imaging Spectrograph (IRIS) and Hinode imaging spectrometer. We summarise how these observations can be coordinated.

Key words. techniques: imaging spectroscopy – Sun: corona – Sun: UV radiation – Sun: heliosphere – instrumentation: high angular resolution – instrumentation: spectrographs

1. Introduction

The Solar Orbiter (SO) space mission (Müller et al. 2020), launched in February 2020, provides both remote sensing and in situ measurements of the solar atmosphere and heliosphere. The main goal of this mission is to understand how the heliosphere is formed and sustained. The trajectory of SO takes it out of the Earth’s orbit and into an orbit around the Sun, reaching to within 0.28 AU of the orbit of Mercury. A spacecraft whose orbit is away from the Sun-Earth line opens up huge possibilities for the representation of the three-dimensional (3D) imaging and spectroscopic data together with near-Earth imaging and spectroscopy. For the first time spectrometers will be located at different angles from each other, allowing 3D spectroscopy of the solar atmosphere.

The presence of persistent high-temperature, high-speed upflows from the edges of active regions (Harra et al. 2008) is a key discovery from Hinode (Culhane et al. 2007). Measurements from the extreme ultraviolet (EUV) Imaging Spectrometer (EIS) indicate that the upflows reach velocities of 50 km s⁻¹ with spectral line asymmetries approaching 100 km s⁻¹ and more (see e.g. Dolla & Zhukov 2011). It has been suggested that these upflows may lie on open magnetic field lines that connect to the heliosphere and may be a significant source of the low-speed solar wind (Harra et al. 2008; Brooks & Warren 2011; Mandrini et al. 2014, and references therein). All of the active regions observed by Hinode/EIS show upflows. Different explanations have been given for the physical mechanism of the upflows, including waves, reconnection in the corona (Baker et al. 2009; Mandrini et al. 2015), and reconnection in the chromosphere (De Pontieu et al. 2007) driving energy upwards. The blue-shifts in these regions are ubiquitous, and indicate the presence of upflows. Various studies have used modelling to determine whether the upflows that are seen at the edges of the active regions become plasma outflowing into the solar wind (e.g. Boutry et al. 2012; Edwards et al. 2016).

Attempts have been made to understand how these flows vary with the location on the disc. Limb-to-limb studies of an active region were carried out by Démoulin et al. (2013) and Baker et al. (2017). The highest plasma velocities in the three spectral lines that they explored have similar magnitudes, and their magnitudes increase with temperature. The authors concluded that their results are compatible with the active-region upflows originating from reconnection between active-region loops and neighbouring loops. However, having two spectroscopic views of an active region will enhance our understanding of flows.

* Movies associated to Fig. 1 are available at https://www.aanda.org

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1,2,3,4,5,6
The first work on solar stereoscopy was carried out by assuming temporal stability of the features under study and allowing the Sun to rotate in order to obtain more than one viewpoint (e.g., Berton & Sakurai 1985; Koutchmy & Molodenskii 1992; Aschwanden & Bastian 1994; Aschwanden et al. 1995, 1999; Feng et al. 2007a). This approach is used also now (see e.g. Nisticò et al. 2013). The review by Aschwanden (2011) presents in detail the methods of static and dynamic solar rotation stereoscopy for coronal loops observed in optically thin coronal plasma EUV emissions.

Two simultaneous viewpoints of the Sun were provided by the Solar Terrestrial Relations Observatory (STEREO, Kaiser et al. 2008). The two STEREO spacecraft orbit the Sun with increasing separation angles, providing stereoscopic images of the Sun’s atmosphere. The stereoscopic images obtained by the two nearly identical EUV broad-band imagers of the Sun Earth Connection Coronal and Heliospheric Imaging System (SECCHI) suite (Howard et al. 2008; Wuelser et al. 2004; Eyles et al. 2009) on board STEREO helped us to understand the 3D geometry of a rich variety of optically thin solar structures (Liewer et al. 2009; Patourakos et al. 2009; Aschwanden 2009a; West et al. 2011; Delanné et al. 2014; Podladchikova et al. 2019; Mierla et al. 2008, 2009, 2010; Temmer et al. 2009; Feng et al. 2009; de Patoul et al. 2013). Methods of stereoscopic triangulation using a stereoscopic pair of EUV or white light coronal images have been developed. A great deal of progress has been made to improve magnetic field models based on stereoscopic information, and to reduce the discrepancy between theoretical magnetic field models and observed stereoscopically triangulated loop 3D coordinates (Aschwanden 2011).

Several studies have been carried out in parallel to calculate the true location of coronal structures in 3D space (Pizzo & Biesecker 2004; Inhester 2006; Feng et al. 2007b; Aschwanden et al. 2008; Howard & Tappin 2008). These studies are based on the direct geometric triangulation using a series of line-of-sight (LOS) measurements taken from different spacecraft views towards the apparent edges of the structures. The true 3D coordinates of the structures are calculated from the intersections of these LOS's. The theoretical background for solar stereoscopy based on the direct triangulation of solar structures is given by Inhester (2006). The direct triangulation technique can be applied after prior identification and matching of the targeted structures in two images (see the review by Wiegelmann et al. 2009). An alternative method of magnetic stereoscopy is proposed by Wiegelmann & Neukirch (2002). The success of magnetic stereoscopy depends on the quality of theoretical magnetic field models. A critical assessment of non-linear force-free field (NLFFF) models has identified a substantial mismatch between theoretical magnetic field models extrapolated from photospheric magnetograms and stereoscopically triangulated loops, of the order of a 3D misalignment angle of $\alpha_{\text{mis}} \approx 20^\circ$ to $40^\circ$ (De Rosa et al. 2009; Sandman et al. 2009; Aschwanden 2011). Rodriguez et al. (2009) highlighted that every pixel in an image is a result of the LOS integration of the emission of the optically thin coronal plasma. This problem of LOS integration persists in the analysis of the EUV data, further complicating the matching of points between the two images.

The triangulation method has been applied to detailed 3D reconstructions of coronal loops in active regions obtained from a stereoscopic pair of images (Inhester 2006; Aschwanden et al. 2008, 2012a, 2015; Aschwanden & Wülser 2011; Nisticò et al. 2013; Chifu et al. 2017). An example of such work (Rodriguez et al. 2009) reveals that loops that appear to be co-spatial in 171 Å and 195 Å images in fact have different heights and occupy different volumes. These results are key to understanding coronal heating.

In this paper we develop a 3D EUV spectroscopy methodology for active regions that is used to build velocity vectors from a pair of EUV images and from Doppler shift maps taken from different perspectives. We do this using the STEREO triangulation technique of EUV images of active regions and we develop novel methods of analytical geometry in space to determine 3D velocities in coronal loops. An advantage of having two different views of an active region is that we can determine the projection angles of the loops, and hence direction of plasma flows along them. We used Hinode/EIS data acquired at different times in order to replicate the situation of two spacecraft with spectrometers on board, and we chose an active region that did not show significant changes in time (Sect. 2). In addition, we describe the methods with which we can reconstruct the velocity vectors; they will be released as a suite of 3D spectroscopy algorithms called the DOPler VElocities Stereoscopically (DOVES) software package, for velocity vector reconstruction. We describe the spatio-temporal co-alignment between EUV broad-band and spectroscopic images (Sect. 3), the reconstruction of the 3D geometry of coronal loops (Sect. 4), and deprojection algorithms of the measured LOS Doppler shifts into velocity vectors of plasma flows onto straight and curved coronal loop structures (Sect. 5). We also describe the magnetic field modelling of the active region to investigate the possibility of using it to derive the 3D loop geometry (Sect. 5.2.4).

A description of the instruments that could be used for 3D stereoscopy is given in Sect. 6. We provide recommendations on the optimum possible spacecraft configuration and spatial resolution of the instruments for 3D spectroscopy within the framework of the Solar Orbiter mission. The proposed 3D spectroscopy method is of particular interest to understanding how the corona is heated and how the solar wind is formed.

2. Data preparation for 3D spectroscopy

2.1. Instrumentation

Initial data for 3D velocity reconstruction includes a stereoscopic pair of EUV images, and simultaneous Doppler shift measurements from two angularly separated spectrometers. To simulate the stereo view we obtain usable simultaneous datasets with SPICE, and we use solar rotation to produce two different viewpoints, as was done for the STEREO software preparation. These viewpoints are imitated now by imaging data obtained by the Atmospheric Imaging Assembly (AIA) on board the Solar Dynamic Observatory (SDO, Pesnell et al. 2012) and spectroscopic data from Hinode/EIS. By applying direct triangulation methods to SDO/AIA intensity images, we restore the 3D coordinates of the observed structures. Then the triangulation of the Doppler shifts observed by Hinode/EIS allows us to restore the vector velocities in these structures.

Hinode is a Japanese mission launched in 2006 by the Institute of Space and Astronautical Science (ISAS) of the Japan Aerospace Exploration Agency (JAXA) in collaboration with the National Astronomical Observatory of Japan (Kosugi et al. 2007). Hinode/EIS has two wavelength bands, 170–211 Å and 246–292 Å, that include spectral lines formed over a wide range of temperatures, from chromospheric to flare temperatures. Hinode/EIS has an effective spatial resolution of about 3–4 arcsec. The high spectral resolution (0.06 Å) allows the determination of Doppler velocity maps with an accuracy of 3 km s$^{-1}$. The Hinode/EIS campaign used in this work has a slit...
Fig. 1. Whole Sun images obtained with SDO/AIA at 193 Å on 19
November 2019 (position 1, left) and 22 November 2017 (position 2, right). The active region (indicated by arrows) has a simple magnetic
configuration ($\alpha$ according to the Zurich classification). These two
views correspond to a separation of 37° between position 1 and posi-
tion 2, allowing the use of stereoscopy for the 3D reconstruction of the
active region. The active region dynamic evolution is illustrated with
four animations: 193 Å from 19 to 22 November (Movie 1), stabilized
193 Å from 19 to 22 November (Movie 2), 2D SDO/HMI photospheric
magnetic field with overlaid 193 Å (Movie 3), 171 Å (Movie 4). The
animations demonstrate that some of the observed simple shape loops
can be traced throughout the entire animation.

size of 1 arcsec, a field of view (FOV) of 467 arcsec $\times$ 511 arcsec
and a raster duration of 67 min.

The SDO/AIA, built by Lockheed Martin Solar and Astro-
physics Laboratory, provides continuous broad-band full-disc
images from the chromosphere to the solar corona in seven EUV
wavelengths covering the temperature range $2 \times 10^5$–$2 \times 10^6$ K, with a cadence of 12 s and a spatial resolution of 1.2 arcsec (cor-
tending to two pixels). The SDO is in a circular geosyn-
chronous orbit at an altitude of 35 800 km.

Active region NOAA AR 2678 was observed simultaneously
by SDO/AIA and Hinode EIS on 19 and 22 November 2017. The
separation of the two observing positions due to solar rotation is
37°, which is within the angular range 9°–120° in which stereo-
scopic vision and 3D spectroscopy becomes applicable. We
discuss this aspect in Sect. 3.

2.2. Observations of NOAA AR 2678

On each viewpoint, three simultaneous images were obtained:
(1) Fe XII intensity from SDO/AIA, (2) Fe XII intensity from
Hinode EIS, and (3) Fe XII Doppler velocity map. The SDO/AIA
images are used to establish the 3D coordinates of coronal points
and the Hinode EIS images for velocity vector construction in
3D space. Figure 1 shows whole Sun images obtained with
SDO/AIA at 193 Å. Because the active region evolves slowly, its
apparent morphological changes are due to the different viewing
angles.

3. Spatio-temporal data co-alignment for 3D spectroscopy

3.1. Hinode/EIS data

We used the eis_prep.pro routine in Solar Software to calibrate the
Hinode/EIS data. The dark current and cosmic rays were
removed, and the hot pixels were corrected. The intensity dig-
ital number (DN) values were calibrated to the spectral radiance
in units of erg (cm$^2$ s sr Å)$^{-1}$. We focused on the strongest emis-
sion line observed by Hinode EIS, Fe XII at 195.12 Å with for-
mation temperatures log $T = 6.1$. For each spectrum we fitted a
single-Gaussian function rather than a double-Gaussian function
with the eis_auto_fit routine and obtained the line peak intensity,
Doppler velocity, and line width using eis_get_fitdata.pro.

3.2. SDO/AIA data processing

We chose the following conditions for the SDO/AIA 193 Å data; the observation time corresponds to that of the whole
Hinode/EIS raster scan. The FOV is slightly larger than the
region observed by Hinode/EIS and it is tracked with solar
rotation at the Carrington rate. Images were transformed to a
grid with a resolution of 0.6 arcsec per pixel. The pre-processed
data were aligned with the solar north and divided by the
exposure time. The strongly saturated frames were removed.
We selected SDO/AIA 193 Å images with more than 2 s of
exposure time to obtain suitable images for long-lived coro-
nal structure analysis. We processed the images with an image
stacking technique to increase photon-to-noise statistics and then
with a wavelet high-pass filter to enhance coronal structures
(Stenborg et al. 2008).

3.3. Hinode/EIS and SDO/AIA data co-alignment

Dynamic structures on the solar surface evolve over a signifi-
cantly shorter time than the duration of the Hinode/EIS raster
observation of 67 min. We created a pseudo SDO/AIA raster
analogous to the Hinode/EIS raster to compensate for possi-
ble dynamic changes. To this end, we extracted the SDO/AIA
data closest in time to each position of the Hinode/EIS slit.
These data were merged into a single SDO/AIA map (hereafter
referred to as the artificial SDO/AIA raster map), corresponding
to the Hinode/EIS raster map using the cross-correlation method
described here. The offset between the artificial SDO/AIA raster
map and the Hinode/EIS raster map was applied, together with
a pointing correction. Based on the new Hinode/EIS pointing
details, the new artificial SDO/AIA raster map was created and
aligned to the new Hinode/EIS map. In our analysis, we
set the expected accuracy to half the size of the Hinode/EIS
pixel.

Figure 2 shows the Hinode/EIS intensity and Doppler veloc-
ity raster maps and the corresponding artificial SDO/AIA raster
map. A similar procedure was presented in Barczynski et al.
(2018) to co-align SDO/AIA and the Interface Region Imaging
Spectrometer (IRIS) raster data. Here 3D point triangulation is
performed with SDO/AIA data, and each 3D point is assigned
to a pair of Doppler shifts via spatio-temporal SDO/AIA–
Hinode/EIS co-alignment.
4. 3D triangulation of NOAA AR 2678

4.1. Triangulation through epipolar geometry

The 3D coordinates are calculated as (Earth-based) Stonyhurst heliographic longitude and latitude, along with the radial distance in solar radii. Then the coordinates are converted into the Heliocentric Earth Equatorial (HEEQ) coordinates for a Cartesian representation of the data. A two-dimensional (2D) solar image taken by a spacecraft is usually identified by the \((i, j)\) coordinates in the image coordinate system or by the latitude and longitude in the heliographic (or helioprojective Cartesian) coordinate system, when \((i, j)\) are projected on the solar surface (Thompson 2006). The multi-point heliospheric observation in the corona requires a complete 3D heliocentric coordinate system, as explained in Sect. 4.1.1. Figure 3 shows the HEEQ coordinate system with the origin in the centre of the Sun and the object position in the heliosphere described by the Cartesian coordinates \(X, Y, Z\). The Stonyhurst coordinate system describes the 3D position of a feature with the spherical coordinates \(R, \Theta, \varphi\).

As Fig. 3 shows, two observing spacecraft positions together with a targeted point in the solar corona define a plane called the epipolar plane. All targeted points have planes in common that contain the two spacecraft positions. Given that every epipolar plane is seen head-on from both spacecraft, it is reduced to a line in the respective image projections. This line is called an epipolar line.

Any targeted corona point found to be situated on a certain epipolar line in one image must lie on the same epipolar line in the other image. The epipolar lines therefore provide a natural coordinate system for stereoscopic reconstructions. Consequently, finding a correspondence between pixels in the images taken by two EUV imagers separated in space is reduced to establishing a correspondence between pixels along the same epipolar lines in the two images.

Once the correspondence between the pixels is found, the 3D reconstruction is performed by calculating the LOSs that belong to the respective pixels in the image and back-tracking them into 3D space. Because the LOSs must lie in the same epipolar plane, their intersection in this plane is defined unambiguously. This procedure is often called ‘tie-pointing’ (see e.g. Inhester 2006; Mierla et al. 2009; Liewer et al. 2009; Aschwanden 2011). The geometrical reconstruction errors are related to both the separation angle \(\gamma\) between the two spacecraft and the spatial resolution \(d_s\) of an image (Inhester 2006; Aschwanden et al. 2015); see Sect. 6.3 for more information on reconstruction errors in the context of the Solar Orbiter mission.

4.1.1. 3D triangulation software

In this study, we use the HEEQ system. The origin of the HEEQ system is the intersection of the solar equator with the central meridian as seen from Earth, and the solar feature location is given either in Cartesian coordinates \(X, Y, Z\) (HEEQ coordinates) or in spherical coordinates (Stonyhurst heliographic coordinates: latitude \(\Theta\), longitude \(\varphi\) and heliospheric radius \(R\) (Thompson 2006) (see Fig. 3). The locations of the features were processed with the World Coordinate System routine `scc_measure.pro` of the Solar Software (Thompson 2006). The structures were traced in stereoscopic pairs of images using projections along the epipolar lines. The routine uses a combination of the information in the header, such as the pixel-to-degree conversion, and the Spacecraft Planet Instrument C-matrix Events orientation (SPICE) database of orbital kernels by calling the routine `convert_sunspice_coord.pro`, which is based on SPICE orbital ephemeris kernels containing the spacecraft location and pointing information. The coordinate system of each spacecraft image plane can be related to a heliocentric coordinate system, and stereoscopic analysis is performed. The output of `scc_measure.pro` is given in the Stonyhurst heliographic coordinates.

4.1.2. 3D triangulation of AR 2678 points

Direct measurement of the 3D coordinates of a point located high in the solar corona is not possible. However, if the two LOSs traced back from the LOS projections on the 2D image intersect higher in the corona at the observed point, then the 3D coordinates of the point can be evaluated. The procedure is performed as follows:

- Point selection on a 2D plane of sky (Image 1); see Fig. 4a.
- Transition to 3D coordinates.
- LOS drawing from satellite 1 to the observed point (LOS 1 or epipolar lines in stereoscopy). The LOSs are not visible on the 2D Image 1 (Fig. 4a), because they are projected to one point.
- Image 2 (Fig. 4b) shows the same area in the solar corona from Perspective 2. LOS1 is visible in Image 2 (Fig. 4b) thanks to observation from a different point.
- The intercept of LOS 1 with coronal loops, the structure of the observation, determines the projection of the 3D point on Image 2 (crosses in Fig. 4b). Thus, the problem converges as long as we deal with two LOS projections of the same point in 3D space.

To check the obtained coordinates the procedure can be then performed in reverse order (Figs. 4c and d). In Fig. 4d we start from the obtained points on Image 2 and describe how we arrive at the initially selected points on Image 1 (Fig. 4c) using the LOS projection technique. We repeat the procedure for points 1, 2, and 3 located on the same loop (and for points 4, 5 and 6, 7) selected on the SDO/AIA images, and express their true spatial position in a 3D heliocentric coordinate system.

4.2. NOAA AR 2678 3D loop triangulation

Here we study the loops, shown in Fig. 4 as straight, curved, or semicircular, which is typical for loops that are inclined towards...
the surface near the centre of the disc (Reale 2010). Aschwanden (2011) presents the technique of semi-circular loop point positions and loop plane inclination definitions. The plane in the strict geometric sense has zero thickness; nevertheless, the terms ‘loop plane’ and ‘loop plane inclination’ are the established and widely used terms in the analysis of the parameters of coronal loops (see e.g. the section ‘Stereoscopic Fitting of Circular Loop Geometry’ in Aschwanden 2009b, or Nisticò et al. 2013; Aschwanden & Wülser 2011), even though the real loops are thick (Klimchuk et al. 1992).

From the 3D triangulation we obtain several points on the same loop. We model the loop as a circle, as explained in Rodriguez et al. (2009) and Aschwanden (2011), defining the curvature radius, and we use three loop points triangulated as described above to fit the loop with the model. The loop can be defined in 3D space by a set of three points with coordinates \((X, Y, Z)\) in the HEEQ coordinate system, and the two sets of 2D coordinates of the solar images \((i_{\text{Los1}}, j_{\text{Los1}})\) and \((i_{\text{Los2}}, j_{\text{Los2}});\) however, it is possible to use more points to fit the circle. We save the projected coordinates of the 3D points \((i_{\text{Los}}, j_{\text{Los}})\) on the images obtained from the two perspectives for the subsequent operations with Doppler velocity measurements. More details on loop parameterisation are given in Aschwanden et al. (2015), Aschwanden (2011), and Nisticò et al. (2013).

When working with optically thin objects in the solar corona, the most critical task is to ensure that two LOSs truly intersect on the object under study in 3D EUV images of the corona. Figure 5 shows the coronal loop under study, presented in the 3D HEEQ coordinate system and reconstructed from points 1, 2, and 3. Blue lines outside the loop show tangents to the loop plotted at the observed points.

Table 1 summarises the results of the 3D reconstructed points in the Stonyhurst coordinate system from the perspective on Earth, and from the perspectives of Position 1 and Position 2. The first entry in the table lists the results for points selected on a closed loop within the active region core, where the dominant Doppler velocity indicates a downflow, and the bottom part of the table shows results for points on open loops at the edges of the active region, where the Doppler velocities show blue-shifted plasma (upflow). We note that we use the term ‘closed loops’ for those for which we can identify their two bases in the AIA EUV images and ‘open loops’ for those for which we can identify only one of their bases in the AIA EUV images.

5. 3D spectroscopy methodology

5.1. Vector velocity measurements

We developed a sequence of algorithms designed to perform triangulation of Doppler shift maps of the solar corona recorded from different perspectives, with the ultimate goal of enabling reconstruction of the vector velocity field.

Solar-rotation stereoscopy was used to reconstruct the 3D coordinates of the highly elevated points in the solar corona. The same loops were identified in two 2D EUV images in order to calculate the true loop location in 3D corona coordinates. We used the coronal loops of NOAA AR 2678 projected on the plane of the sky observed by SDO/AIA from two different viewing angles (Figs. 6a and b). The region was stable in time and did not show any flaring activity, as demonstrated by Animation 1. The crosses in the images show the different locations of the elevated coronal points when they are projected through the LOS onto 2D images. The SDO/AIA images and Hinode/EIS intensity images were co-aligned as described in Sect. 3 (Figs. 6c and d). As long as the Hinode/EIS intensity and velocity maps are co-aligned, the crosses on the intensity maps can be directly translated onto the velocity maps. Thus, the crosses on the two Hinode EIS velocity maps indicate two LOS velocities measured in the considered elevated coronal points (Figs. 6e and f).

5.2. Velocity vectors for NOAA AR 2678 coronal loops

In this section we describe how to obtain a velocity vector at a point situated high above the photospheric level, using two LOS

Fig. 4. (a) Points 1 to 7 on Image 1 selected for triangulation. (b) Points 1 to 7 on Image 2 are found on the corresponding epipolar lines (dashed lines) at the intersection with the observed structures. Points 1, 2, and 3 belong to closed loops (i.e. where both of their bases can be identified in AIA EUV images); whereas points 4, 5, 6, and 7 belong to open loops (i.e. where only one of their bases can be identified AIA EUV images). To validate the tie-pointing procedure the reverse process was performed, first starting from the obtained points in Image 2 (panel d) and then selecting their stereoscopic pairs on the corresponding epipolar lines in Image 1 (panel c), obtaining the same points.

Fig. 5. Semi-circular curved coronal loop under study observed stereoscopically in NOAA AR 2678 in November 2017. The loop is shown in the HEEQ coordinate system. The blue lines outside the loop show tangents to the loop plotted at the selected points.
Table 1. NOAA AR 2678 3D loop triangulation results.

| Closed loop | Position 1 (Lat./Long.) | Position 2 (Lat./Long.) | Height above photosphere | Inclination α degrees |
|-------------|------------------------|------------------------|-------------------------|----------------------|
| Point 1     | $-10.7°/-3.04°$        | $-13.2°/-3.13°$        | $12.7°/35.16°$          | 16.0 Mm              |
| Point 2     | $-12.9°/5.2°$          | $-16.6°/-5.9°$         | $16.2°/36.40°$          | 76.1 Mm              |
| Point 3     | $-12.9°/-12.2°$        | $-15.6°/-12.9°$        | $15.3°/26.2°$           | 29.6 Mm              |

Notes. Latitude, longitude, and height above the photosphere of the seven selected points in Fig. 4. Column 1 indicates the point number. Columns 2, 3, and 4 show the latitude and longitude from Earth, Position 1 viewpoint, and Position 2 viewpoint, respectively. Column 5 lists the height of the points above the photosphere, and Col. 6 shows the inclination angles $\alpha$ between either the open or closed coronal loop plane and the vertical to the solar surface. See text for the meaning of closed and open loops.

Doppler measurements obtained by spectrometers distributed in the heliosphere. Previously, we restored the true 3D coordinates of points in the high corona and the coordinates of their projections onto each of the 2D Doppler maps.

5.2.1. Deprojection rules of two LOS velocity components in the epipolar plane

Observing two LOS Doppler velocities at any identifiable point in the corona enables us, first of all, to find the velocity vector in the epipolar plane formed in the 3D heliosphere by the position of the two satellites and the observation point. The grey plane in Fig. 7 shows the epipolar plane defined by the observed point 4 on the open coronal loop and the location of the two satellites. The two Doppler blue-shifted velocities measured at point 4 along the LOSs situated between the observed point and the observing satellites represent the 1D vectors $v_{LOS1}$ and $v_{LOS2}$. Figure 8 shows two LOS Doppler velocities measured at points 4, 5, 6, and 7. Here, $v_{LOS1}$ and $v_{LOS2}$ are different velocities measured along the LOSs. The velocity vector $v_{ep}$ is the projection of the 3D velocity in the epipolar plane, and can be reconstructed by calculating the intersection point of the two perpendiculars deprojected from $v_{LOS1}$ and $v_{LOS2}$ in the epipolar plane. The velocity vector $v_{ep}$, defined in the epipolar plane, is an especially useful estimate when we cannot identify any structure to which this point belongs. In this case $v_{ep}$ can be an indicator of the plasma flow behaviour.
5.2.2. Plasma outflows along open loops

In this subsection we demonstrate the technique of velocity vector reconstruction in the open coronal loops using the pre-calculated geometry in 3D space. If the pair of points (4, 5) or (6, 7) belongs to the open coronal loop approximated as a straight line between two points, we can restore the velocity vector magnitude of the plasma flow by the direct deprojection of \( \mathbf{v}_{ep} \) on the loop by using the formula

\[
\mathbf{V}_{3D} = \frac{\mathbf{v}_{ep}}{\cos(\beta)},
\]

where \( \beta \) is the angle between the \( \mathbf{v}_{ep} \) velocity and the coronal loop (Fig. 9). If \( \beta \approx 90^\circ \), the deprojection cannot be made. Correspondingly, small \( \beta \) angles lead to higher accuracy of the \( \mathbf{V}_{3D} \) deprojection.

We reconstructed the velocity vector \( \mathbf{V}_{3D} \) for the points in the open-loop outflows using deprojections: \( \mathbf{v}_{ep} \rightarrow \mathbf{V}_{3D} \). Figure 9 shows the locations of the coronal loop in 3D space and the epipolar plane (grey). The epipolar plane is defined by the locations of the two spacecraft and the point in the solar corona to be triangulated. \( \mathbf{v}_{ep} \) represents the projection of \( \mathbf{V}_{3D} \) on the epipolar plane. When projected on LOS1 and LOS2, \( \mathbf{v}_{ep} \) represents the velocities \( \mathbf{v}_{LOS1} \) and \( \mathbf{v}_{LOS2} \) measured by each spectrometer.

5.2.3. Plasma outflows in closed loops

In this subsection we demonstrate a technique for determining the velocity vector for flows following curved paths traceable in the solar corona. When measuring the LOS Doppler velocities with two spectrometers at points belonging to a closed, semi-circular, or simply curved coronal structure, \( \mathbf{V}_{3D} \) can be determined again using Eq. (1). However, in this case the deprojection should be performed on the tangent to the coronal-loop direction, as long as the velocity vector is oriented in this direction. Figure 10 (left panel) shows the deprojected velocity \( \mathbf{v}_{ep} \) from the two measured LOS Doppler velocities \( \mathbf{v}_{LOS1}, \mathbf{v}_{LOS2} \). Figure 10 (right panel) shows \( \mathbf{V}_{3D} \) at points 1, 2, and 3. Detailed information on the LOS Doppler velocity magnitudes and their angles with the tangents are listed in Table 3.

Thus, the velocities deprojected from two LOS Doppler velocities measured at points of a closed coronal loop provide true intrinsic information of the plasma flow behaviour in the loop. In this case the deprojection is not carried out on the structure, as described in the previous subsection, but in the direction of the tangent to the loop reconstructed at the observed point. This technique can be applied to any curved structure in the solar atmosphere, provided that we can define its 3D geometry using, for instance, epipolar geometry stereoscopy or magnetic modelling of the loops.

5.2.4. Comparison of deduced inclinations and linear force-free field extrapolation of NOAA AR 2678

Magnetic field modelling has been used extensively to interpret the morphology of coronal magnetic loops (Chiu et al. 2015). An example by Harra et al. (2008) illustrates how using the angles from magnetic field modelling enables a more accurate determination of the upflowing plasma at the edge of an active region.

We compute the coronal magnetic field topology of NOAA AR 2678 during its disc transit on 11 November 2017 at 08:28 UT and 22 November 2017 at 03:44 UT. The LOS magnetic field is extrapolated to the corona using a linear force-free field (LFFF) configuration where \( \mathbf{J} \times \mathbf{B} = 0 \) and \( \nabla \times \mathbf{B} = \alpha \mathbf{M} \), with \( \alpha \) constant (Mandrini et al. 1996; Démoulin et al. 1997).

Figure 11 (left panel), shows the LFFF model results for 19 November 2017. Similar results are obtained on 22 November 2017 (right panel). Both panels show the modelled global coronal structure for which we have used the magnetograms closest in time to each EIS image as boundary condition. The panels are constructed superposing field lines computed using different values of \( \alpha \), which have been selected to better match the shape of the observed loops in AIA 193 Å. To do this comparison, the model is first transformed from the local frame to the observed frame, as discussed in Mandrini et al. (2015) (see the transformation equations in the appendix of Démoulin et al. 1997). This allows a direct comparison of our computed coronal field configuration to AIA EUV loops. Furthermore, in order to determine the best matching \( \alpha \) values we have followed the procedure discussed by Green (2002).

Figure 12 shows the definition of the loop inclination angles \( \delta_{\text{EUV}} \) computed for EUV SDO/AIA loops using triangulation methods, which are compared with those \( \delta_{\text{M}} \) derived from our LFFF extrapolation. The inclination angles \( \delta_{\text{M}} \) derived from the model are computed as discussed in Démoulin et al. (2013); they are the angles between the line tangential to the loop at the coronal points studied here and the vertical to the solar surface as shown in Fig. 12.

The LFFF and stereoscopic methods both obtained rather close ranges of inclination angles to the vertical: [20°, 30°] at point 1 and [35°, 45°] at point 3 on the closed loop. For open loops the inclination \( \delta_{\text{EUV}} \) and \( \delta_{\text{M}} \) to the vertical at point 6 and point 7 (where the field lines are quite simple) are in the same range [10°, 30°]. At point 4 and point 5 the inclination angles...
estimated by the two methods are in the range [15°, 25°]. Since the projection of an observed coronal loop has a certain thickness, a computed set of infinitely thin field lines can match the visible shape of a loop; this is the origin of the range of values for the inclination \( \delta_M \) coming from the magnetic field model.

In general, we do not observe discrepancies between the field lines computed from the LFJFF model and stereoscopic reconstruction, although photospheric footpoints are difficult to identify in EUV images, and from the magnetic model we trace the full line starting at the photospheric level. The latter demonstrates the future opportunity of pairing stereoscopic loop reconstruction using Solar Orbiter data with magnetic field modelling, which is expected to significantly improve the accuracy of stereoscopic vector velocity measurements (see e.g. Aschwanden et al. 2015).

6. 3D Spectroscopy by two distributed imaging spectrometers

We developed an approach using a pair of space-born spectrometers that provide stereoscopic views of Doppler shift maps in the solar corona. Measurement of the velocity vector in the corona can be achieved with one extraterrestrial remote spectrometer (SO/SPICE) and with the orbital spectrometer Hinode/EIS or IRIS. Dual spectrometers allow measurement of the 2D velocity vector in any identifiable coronal point, and eventually reconstruction of 3D velocity vectors, using the morphology of the coronal structures. In this section we describe how we derive the optimal spacecraft configurations for 3D spectroscopy with SO.

The spectral and spatial resolution of spectrometers influences directly the velocity vector measurement accuracy.

6.1. The SPICE, Hinode/EIS, and IRIS spectrometers

High-resolution spectroscopy of optically thin plasma allows measurements of plasma parameters such as temperatures, densities, chemical abundances, and Doppler and non-thermal motions (Phillips et al. 2008; Del Zanna & Mason 2018).

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**Table 2. Plasma flow velocities reconstructed in NOAA AR open loop 3D points.**

| Open loop | \( V_{LOS1} \) (km s\(^{-1}\)) | \( V_{LOS2} \) (km s\(^{-1}\)) | \( v_{ep} \) (km s\(^{-1}\)) | \( \beta \) (degrees) | \( V_{3D} \) (km s\(^{-1}\)) |
|-----------|-------------------------------|-------------------------------|----------------------------|-----------------|-----------------|
| Point 4   | -13.61 km s\(^{-1}\)         | -16.71 km s\(^{-1}\)         | -16.73 km s\(^{-1}\)      | 28.86\(^\circ\) | -19.11 km s\(^{-1}\) |
| Point 5   | -4.54 km s\(^{-1}\)          | -5.67 km s\(^{-1}\)          | -5.67 km s\(^{-1}\)       | 28.54\(^\circ\) | -6.45 km s\(^{-1}\) |
| Open loop | \( b \)                       | \( b \)                       | \( b \)                     | \( b \)         | \( b \)         |
| Point 6   | -3.37 km s\(^{-1}\)          | -5.27 km s\(^{-1}\)          | -5.45 km s\(^{-1}\)       | 36.29\(^\circ\) | -6.76 km s\(^{-1}\) |
| Point 7   | -2.03 km s\(^{-1}\)          | -3.17 km s\(^{-1}\)          | -3.29 km s\(^{-1}\)       | 36.32\(^\circ\) | -4.07 km s\(^{-1}\) |

**Notes.** Numerical values of the fully reconstructed velocity vectors \( V_{3D} \), which characterise outflows along open loops shown by black arrows in Fig. 9, \( \beta \) is the angle between the velocity vector \( v_{ep} \) measured in the epipolar plane and the loop. \( V_{3D} \) can be reconstructed by deprojecting \( v_{ep} \) onto the coronal loop. The calculations were carried out for each of the points 4, 5, 6, and 7, which lie on open loops. The minus sign is implied to indicate outflows.
In the science phase of the Solar Orbiter mission, there will be opportunities to carry out 3D spectroscopy in the chromosphere, transition region, and corona for the first time by combining data from Hinode/EIS (Culhane et al. 2007), IRIS (De Pontieu et al. 2014a), and SO/SPICE (Anderson et al. 2020). The Hinode/EIS wavelength ranges are 1170–210 Å and 250–290 Å, and its angular resolution is 2 arcsec. There are four slit and slot positions: 1 arcsec slit, 2 arcsec slit, 40 arcsec slot, and 266 arcsec slot. The temporal resolution is a few seconds in dynamic events, 10 s in active regions, and around 1 min in coronal holes. The maximum FOV is 360 arcsec × 512 arcsec.

The SO/SPICE instrument is also a high-resolution imaging spectrometer operating at EUV/UV wavelengths. Its design is focused on studies that combine remote sensing and in situ instruments on board Solar Orbiter. The wavelength range is 704–790 Å and 973–1049 Å and the accuracy of the line shifts is 5 km s\(^{-1}\). Table 4 shows the different emission lines used by SPICE, their wavelengths, temperature of formation, and their closest equivalents for Hinode/EIS. In most cases the emission lines for the two instruments do not represent the same ion, but the formation temperature is similar. Some caution must be used for lines that are optically thick, such as He II, when comparing them.

IRIS is an imaging spectrometer with a slit-jaw imaging system that probes the chromosphere. It explores the solar chromospheric dynamics to determine how the energy flows through the chromosphere and the transition region. The spacecraft has a polar sun-synchronous orbit. IRIS has a spatial resolution of 0.33–0.4 arcsec, a temporal resolution of 2 s, and a velocity resolution of 1 km s\(^{-1}\) over an FOV of up to 175 arcsec × 175 arcsec (De Pontieu et al. 2007, 2014b).

IRIS and Hinode/EIS frequently carry out simultaneous observations, pointing at the same target. This mode of collaboration will be of great interest alongside the different FOVs of Solar Orbiter. The IRIS wavelength range covers the chromosphere and the Hinode/EIS wavelength range covers the corona. The SPICE wavelength range provides measurements of the chromosphere, transition region, and corona. Combining the three spectrometers will provide the first 3D solar spectroscopy.

Table 3. Plasma flow velocities reconstructed in NOAA AR closed loop 3D points.

|         | \(v_{z,los1}\) (km s\(^{-1}\)) | \(v_{z,los2}\) (km s\(^{-1}\)) | \(v_{\beta}\) (km s\(^{-1}\)) | \(V_{3D}\) (km s\(^{-1}\)) |
|---------|-------------------------------|-------------------------------|----------------------------|----------------------------|
| Closed loop |                               |                               |                           |                           |
| Point 1  | 2.38 km s\(^{-1}\)          | -0.68 km s\(^{-1}\)          | 4.91 km s\(^{-1}\)      | 7.51°                      | 9.4 km s\(^{-1}\)       |
| Point 2  | 0.71 km s\(^{-1}\)          | -1.97 km s\(^{-1}\)          | 4.27 km s\(^{-1}\)      | 17.98°                     | 4.8 km s\(^{-1}\)       |
| Point 3  | -7.35 km s\(^{-1}\)         | -9.4 km s\(^{-1}\)          | -9.4 km s\(^{-1}\)      | 14.11°                     | -9.70 km s\(^{-1}\)     |

Notes. Velocity vectors characterising the flows in the investigated closed loops observed in NOAA AR 2678 in November 2017. \(V_{3D}\), shown in Fig. 10, is reconstructed at points 1, 2, and 3 on the loop. The velocity vectors \(V_{3D}\) are found by deprojection of \(v_{\beta}\) on the tangent to the coronal loop directions. The minus sign indicates outflows.
to the low corona using three imaging telescopes: the Full Sun Imager (FSI) and two high-resolution imagers (HRI\textsubscript{EUV} and HRI\textsubscript{Ly\alpha}). The FSI provides a 3.8° × 3.8° FOV in the 174 Å and 304 Å passbands with a typical temporal cadence of 600 s and a spatial resolution of 0.5 arcsec. This telescope provides a context view of the upper solar atmosphere up to 4 \( R_\odot \) at perihelion.

Both HRI\textsubscript{EUV} and HRI\textsubscript{Ly\alpha} telescopes provide a 17 arcmin × 17 arcmin FOV with a plate-scale of 0.5 arcsec and unprecedented imaging cadence of 1 s. The passband of HRI\textsubscript{EUV} (174 Å) is comparable to that of FSI. The second high-resolution telescope, HRI\textsubscript{Ly\alpha}, observes the Sun in the Ly\alpha line (Ly\alpha 1216 Å). The HRI instruments will study the small-scale structures and highly dynamic events in the upper solar atmosphere.

Ly\alpha is the most intense emission line in the solar spectrum and affects planetary atmospheres, so it is of great interest; however, to date, observations in Ly\alpha are rare and were carried out mainly aboard sounding rockets (Korendyke et al. 2001; Vourlidas et al. 2010, 2016; Chua et al. 2013; Chintzoglou et al. 2018) and the Transition Region and Coronal Explorer space observatory (Golub et al. 1999). Table 6 shows the closest corresponding SPICE emission line to the EUI passband. Again, caution must be used for the temperature of formation with Ly\alpha, as it is a complex, optically thick spectral line.

### 6.3. Spacecraft separation and spatial resolution

#### 6.3.1. 3D point triangulation uncertainties

An issue of great importance when planning observations with two spacecraft is understanding the optimum separation between them. In this section we discuss how the accuracy of the measurement varies with spacecraft separation. Inhester (2006) calculated that the errors in 3D coronal point triangulation depend on both the spatial resolution \( ds \) of an image and the spacecraft separation angle \( \gamma \):

\[
\varepsilon_\Delta = \frac{ds}{\sin(\gamma/2)}.
\]

Here \( ds \) is the pointing error across the epipolar line, which is linearly dependent on the spatial resolution of the instrument and can have values from 1 to 10 pixels, depending on the accuracy with which the same feature in two images can be identified. When the spatial resolution is fixed and \( \gamma = 0 \), the errors are very large. For large \( \gamma \) the errors are small, but additional errors are introduced due to complications in identifying the same feature in both images. When \( \gamma = 180^\circ \), the errors are very large for off-limb features because the LOSs of the two spacecraft are parallel. The low optical thickness of the solar corona enables the observation of the same off-limb feature when the separation is near 180°, but makes it impossible to observe the same on-disc feature (Aschwanden et al. 2012b, 2015).

### 6.3.2. 3D coronal Loop triangulation uncertainties

The geometric error in Eq. (2) is the error for the triangulation of a single-point object. For a 1D object such as a section of a loop, the error is more complex, as it additionally increases with decreasing angle between the loop tangent and the epipolar plane, and theoretically becomes infinite if this angle becomes zero. The intuitive reason for this behaviour is that a smaller angle with the epipolar plane results in a smaller visual angle on the epipolar line, and theoretically becomes infinite if this angle becomes zero. By combining the epipolar line, which is disturbed by projection effects and uncertainty in the identification of stereoscopic correspondence, and the spatial resolution of the instrument and can have values from 1 to 10 pixels, depending on the accuracy with which the same feature in two images can be identified. When the spatial resolution is fixed and \( \gamma = 0 \), the errors are very large. For large \( \gamma \) the errors are small, but additional errors are introduced due to complications in identifying the same feature in both images. When \( \gamma = 180^\circ \), the errors are very large for off-limb features because the LOSs of the two spacecraft are parallel. The low optical thickness of the solar corona enables the observation of the same off-limb feature when the separation is near 180°, but makes it impossible to observe the same on-disc feature (Aschwanden et al. 2012b, 2015).

We used Eq. (2) to calculate the triangulation errors for different spacecraft separation angles \( \gamma \) between the Solar Orbiter and the near-Earth spacecraft Hinode or IRIS. Figure 13 shows the triangulation error as a function of spacecraft separation angle, where \( ds \) is 1 pixel. The physical measure contained in a one-pixel unit will vary as (100–350) \( km^2 \) for HRI and as (900–3000) \( km^2 \) for FSI when the Solar Orbiter travels between aphelion and perihelion. The angle interval of 10°–100° is expected...

### Table 4. SPICE emission lines, wavelength, temperature of formation (T), and the closest equivalents in Hinode/EIS.

| SPICE | Wavelength (Å) | log T | EIS | Wavelength (Å) | log T |
|-------|---------------|-------|-----|---------------|-------|
| C\textsc{iii} | 977.03 | 4.5 | He\textsc{ii} | 256.7 | 4.7 |
| O\textsc{v} | 760.43 | 5.4 | O\textsc{v} | 192.9 | 5.4 |
| Ne\textsc{viii} | 770.92 | 5.8 | Mg\textsc{vii} | 278.39, 280.75 | 5.8 |
| Si\textsc{vii} | 275.35 | 5.8 | Fe\textsc{viii} | 185.21, 194.66 | 5.8 |
| Mg\textsc{ix} | 706.02 | 6.0 | Fe\textsc{x} | 184.5, 257.2 | 6.0 |
| Fe\textsc{x} | 1028.04 | 6.0 | | | |
| Mg\textsc{xii} | 997.44 | 6.2 | Fe\textsc{xiii} | 196.54, 202.04, 203.83 | 6.2 |
| Si\textsc{xii} | 520.67 | 6.3 | Fe\textsc{xiv} | 274.20, 264.79 | 6.2 |
| Fe\textsc{feiviii} | 974.84 | 6.9 | Ca\textsc{xvii} | 192.82 | 6.7 |
| Fe\textsc{xx} | 721.55 | 7.0 | Fe\textsc{xix} | 263.76 | 7.2 |
| | | | Fe\textsc{xxiv} | 255.11 | 7.2 |

### Table 6. SPICE emission lines, wavelength, and temperature of formation, and the closest equivalent filter in EUI.

| SPICE | Wavelength (Å) | log T | EUI | Wavelength (Å) | log T |
|-------|---------------|-------|-----|---------------|-------|
| C\textsc{iii} | 977.03 | 4.5 | Ly\alpha | 1216 | 4.7 |
| Mg\textsc{viii} | 772.31 | 5.9 | Fe\textsc{ix} | 174 | 5.9 |
| Ne\textsc{viii} | 772.31 | 5.9 | Fe\textsc{ix} | 174 | 5.9 |

...
to provide a sufficiently small triangulation error and good pointing possibilities for spectrometers.

7. Discussion and conclusions

We have developed a technique that reconstructs velocity vectors in coronal loops from the Doppler maps and EUV images obtained by two spacecraft located at significant distances from each other (on a heliospheric scale). The methodology is built upon the dynamic solar rotation EUV spectroscopy and on STEREO triangulation methods for coronal loops observed in optically thin coronal EUV emissions lines (Aschwanden 2011). Novel algorithms use methods of spatial analytical geometry in order to determine 3D velocities in coronal loops using different LOSs Doppler shifts measurements. The study is inspired by the launch of the Solar Orbiter in February 2020, the first solar mission to travel far beyond the ecliptic plane, which carries the separation angles between two spacecraft and is therefore not recommended. However, the angle interval of 10°−100° provides a sufficiently small triangulation error and good pointing possibilities for spectrometers.

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