Multimaterial Microfluidic 3D Printing of Textured Composites with Liquid Inclusions

Xiyi Li, Jia Ming Zhang,* Xin Yi,* Zhongyi Huang, Pengyu Lv, and Huiling Duan

Natural materials with heterogeneous and hierarchical structures exhibit outstanding mechanical properties, including high specific stiffness, strength, and damage tolerance. For example, dactyl clubs of harlequin mantis shrimps have significantly high specific strength and toughness due to the inorganic ion contents and spatial variation of the mineral phase in the organic matrix, and the graded and layered structures of the dactyl clubs could retard crack propagation. Though many fabrication methods have been proposed to produce biomimetic engineered materials, composite materials with finely controlled microstructures far exceed the fabrication ability of the conventional manufacturing methods.

Recent developments on the structure design and multiple material printing in additive manufacturing indicate that direct ink writing, an extrusion-based additive manufacturing technique, can potentially serve as a promising way to fabricate composites with finely controlled microstructures. It has been reported that the direct ink writing utilizing inks of different mechanical and physical properties has been used to fabricate particle laden functional porous materials, flexible microelectrodes, microcapacitors, electromagnetic structures, and 3D gradient materials. Solid inclusions such as carbon fibers, lead-free nanowires, and liquid crystals dispersed in the inks can be reorientated along the printing direction under the shear and extensional flow in the printhead nozzle. The spatial orientation of the carbon fibers in the epoxy matrix can also be controlled locally by rotating the nozzle during printing. Moreover, for functionalized fibers, external force fields can be used to control their orientation in the inks. For example, the ultrasonic standing wave is used to control the arrangement of glass fibers in photocurable resin. Magnetized stiff platelets are orientated in the printed objects by applying low magnetic fields. Orientational control over carbon nanotubes in the electrical field is used to fabricate biomimetic architectures. Despite recent advances on the local control over fiber orientation in additive manufacturing, multimaterial 3D printing with a high degree of spatial and compositional precision has not been accomplished, which could open new opportunities for fabricating composite systems of excellent functional properties.

Taking advantage of the fine manipulation of droplets in microfluidics, here we combine direct ink writing and microfluidics to build a multimaterial 3D printing system for fabricating textured composites with liquid inclusions of programmable spatial distribution and compositions. Depending on the geometry of the nozzle, flow rate, printing speed, and spatial selection, structures with different inclusion packing patterns are printed. A combination of simple theoretical analysis and experimental studies is employed to determine quantitative phase diagrams serving as guide in the rational selection of desired printing parameters. Moreover, 1D, 2D, and 3D structures programmed with inclusion packing patterns and compositions are fabricated. Finally, we demonstrate new implications of this microfluidic-based versatile printing framework.
in fabricating layered composite beams of tunable thermal expansion property and self-healing materials. As the liquid inclusions are generated in the form of droplets in the microfluidic part of our printing system, we use droplets and liquid inclusions interchangeably hereinafter except in the discussion on the contact angle measurement of the matrix material.

**Figure 1** schematically illustrates the multimaterial 3D printing system consisting of a microfluidic part, a 3D printing part, and a custom-designed single chip microcomputer (SCM). In the microfluidic part, the inner and outer fluids are injected independently at controlled flow rates by syringe pumps (LSP02-2A, Longer Precision Pump) into the inlet channels and meet at the T-junction in a generator of on-demand mixing droplets which is integrated to produce liquid inclusions of simple and mixed compositions in printed structures (see Figure S1, Supporting Information). In the droplet generator, two inner phases (red and blue) are mixed first to produce (purple) mixed output which is then sheared into droplets by the continuous immiscible outer phase at the T-junction. Depending on the flow rate ratio between the red and blue inner phases, the output droplets can change their color (indicating the degree of mixing) from red to purple, then to blue in a discontinuous 1D line, the nozzle motion along X- and Z-axes. Arrows in the microfluidic part represent the flow direction.

two-part curing agents are encapsulated separately in two different kinds of liquid inclusions and enable a re bonding of the crack face with their release in the damage region.

Water-in-resin two phase fluids are exemplified as the ink in our printing demonstrations, where the inner phase is the water–glycerol mixture and the outer phase is the commercial photosensitive liquid resin which can be regarded as a Newtonian fluid before it is cured (Figure S2, Supporting Information). In the dripping regime, the droplet diameter $D_{\text{drop}}$ scales as $D_{\text{drop}} \approx w_{\text{drop}}^{-1/3}$ (Figure S3, Supporting Information), where $w_{\text{drop}}$ is the outlet channel width and the outer phase capillary number $Ca_{o} = \eta_{o} \mu_{o} / \gamma$ denotes the relative strength between the viscous force and interfacial tension. Here $\eta_{o} (\approx 0.53 \text{ Pa s})$ and $\mu_{o}$ represent the outer phase viscosity and flow velocity at the T-junction, respectively, and $\gamma (\approx 0.02 \text{ N m}^{-1})$ is the interfacial tension between the inner and outer phases. In our system, the outlet channel width is $w_{o} = 330 \mu\text{m}$, and the diameter of the embedded droplets ranges from 200 to 500 $\mu\text{m}$ (Figure S3, Supporting Information). During the printing process, droplets in the gentle flow ($Ca_{o} < 0.1$) could pass through the printhead nozzle without breaking due to a high degree of the droplet deformability.

We first benchmark printing on the glass substrate 1D composite lines and 2D carpet-like structures consisting of water droplets embedded in the resin matrix. During the printing of 1D lines, the nozzle of a diameter $D_{\text{nozzle}} = 0.33 \text{ mm}$ moves at a tip height $H$ and a printing speed $v$. The contact angle of resin on the glass substrate is $38^{\circ} \pm 1^{\circ}$ at $25^{\circ} \text{C}$. As the height of the printed line is less than the characteristic capillary length (1.4 mm), the gravity effects can be ignored and the printed line adopts a cylindrical cap shape of width $w$, contact angle $\theta$, and height $h = w (1 - \cos \theta) / \sin 2 \theta$, and its cross-sectional area is $A = w^{2} (\theta / \sin \theta - \cot \theta) / 4$. At a total flow rate $Q = Q_{i} + Q_{o}$, $A = Q / v$. Therefore, we have $w = 2 \sqrt{Q / v \theta / \sin \theta - \cot \theta}$. This linear relationship between $w$ and $\sqrt{Q / v}$ is confirmed by the experimental results (Figure S4a, Supporting Information). At a large tip height $H > h$, the extruded ink could form a droplet hanging at the nozzle tip and lose contact with the printed line segment as the nozzle moves. Therefore, a discontinuous 1D line with multiple segments (regime I in Figure 2a) is printed as $h < H$, and the corresponding criterion is $H > w (1 - \cos \theta) / (2 \sin \theta)$ or

$$H \geq \frac{1 - \cos \theta}{\sqrt{2} \theta - \sin \theta} \frac{\pi c}{4v},$$

where $c = Q / (\pi D_{\text{nozzle}}^{2} / 4)$ is the extrusion speed of the composite ink from the nozzle.

At $D_{\text{drop}} < w$, we have $V_{\text{droplet}} / V_{\text{resin}} = Q / Q_{o}$. As $V_{\text{resin}} = A \cdot \Delta x - V_{\text{droplet}}$, the droplet spacing $\Delta x$ along the printed line is $\Delta x = (\pi D_{\text{drop}}^{2} v / 6Q_{o})$, consistent with the experimental results (Figure S4b, Supporting Information). In the case of $\Delta x < D_{\text{drop}}$,}

![Figure 1](image-url)
the droplets contact or overlap and form an irregular packing configuration with merged inclusions in the printed line. The criterion for the irregular inclusion packing (regime III in Figure 2a) is

\[ \Delta x < (\pi D_{\text{drop}}^2 v) / (6 \sigma_{\text{air}}) < D_{\text{drop}} \] or

\[ \frac{v}{c_i} < \frac{3 D_{\text{nozzle}}^2}{2 D_{\text{drop}}} \]  

(2)

where \( c_i = Q_i / (\pi D_{\text{nozzle}}^2 / 4) \) is the extrusion speed of the inner phase.

In the rest regime of \( \Delta x > D_{\text{drop}} \) and \( H < h \) (regime II in Figure 2a), the line with regularly packed liquid inclusions is printed as theoretically predicted.

In the fabrication of a 2D carpet-like structures consisting of liquid inclusions encapsulated in the resin matrix, we program parallel straight paths of the nozzle motion at \( H = 150 \, \mu\text{m} \) and conduct printing in the regime II in Figure 2a (regular packing regime). In addition to the parameters of line width \( w \) and droplet diameter \( D_{\text{drop}} \) as discussed in 1D line printing, a new length parameter, the distance \( \Delta y \) between two adjacent printing paths, emerges and regulates the printing of 2D structures. In the cases of \( \Delta y > w \), the printed lines do not overlap and stand alone. In these circumstances, the droplet of diameter \( D_{\text{drop}} > w \) would form individual clusters and a discontinuous line is printed along each printing path (regime I in Figure 2b); for the droplet of diameter \( D_{\text{drop}} < w \), the droplets are regularly packed in the printed lines standing alone (unconnected regular packing regime II arising, see Figure 2b). Note that the droplet inclusions in the discontinuous line in Figure 2a remain intact, while these in discontinuous line in Figure 2b break up.

In the cases of \( \Delta y < w \), the printed lines in adjacent printing paths would form contact with each other and a resin meniscus forms attaching to the nozzle tip and moves during the printing process. Experimental observation indicates that the characteristic size of the meniscus is around \( l \approx 0.18w \). In these cases, droplets of \( D_{\text{drop}} > \Delta y \) might overlap and an irregular packing arises; at \( D_{\text{drop}} < \Delta y < D_{\text{drop}} + l \), the moving meniscus near the nozzle induces an irregular motion of the existing droplets and an irregular packing pattern forms. In both cases, we can observe the regime IV in Figure 2b.

In the rest regime of the \( D_{\text{drop}} / \Delta y - w / \Delta y \) space (\( w / \Delta y > 1 \) and \( D_{\text{drop}} / \Delta y < 1 - 0.18w / \Delta y \)), the printed lines contact and a 2D carpet-like structure is formed with droplets dispersed in a regular pattern (connected regular packing, regime III in Figure 2b).

Based on the analyses above, a packing phase diagram in 2D printing is determined and confirmed experimentally (Figure 2b) and the corresponding printed structures can be found in Figure 2c. The benchmark printing above with corresponding simple analyses show that the liquid inclusion arrangement in the printed objects could be precisely programmed. To validate this in situ control, we print a structure with programmed droplet size and spacing by simply changing the inner phase flow rate (Figure 2d). Decreasing the inner phase flow rate results in a decreased droplet size and an increased droplet spacing at a given printing speed.
In the above experiments, the density of the droplets is 1.138 kg cm\(^{-3}\), and the resin matrix has a density of 0.98 kg cm\(^{-3}\) and a viscosity of 530 mPa s at 26 °C. Due to the small density difference between the droplets and resin matrix and the high viscosity of the resin matrix, the embedded droplets are stable and could remain intact even as the printed line is lifted and deformed (Figure S5b, Supporting Information). In the case where the density difference between the inclusions and the fluidic matrix is significantly different, the inclusions of a smaller density rise to the matrix surface while the inclusions of a larger density could sink. As indicated in our additional experiment (Figure S5c, Supporting Information), the bubble inclusions initially in the resin matrix rise due to their much smaller density and then break at the resin surface.

The in situ programmable generation of liquid inclusions in the matrix material offers the ability to print more complex structures with locally tunable droplet composition and distribution. To demonstrate these abilities, we print 1D, 2D, and 3D objects containing liquid inclusions. In Figure 3a and Video S1 (Supporting Information), we fabricate a spiral structure with droplet inclusions evenly spaced along the structure curve. In a programmed continuous printing, we fabricate a 2D structure with droplet inclusions in a pattern showing the letters PKU (Figure 3b and Video S2, Supporting Information) and with droplets of colors gradually varying from red to blue (Figure 3c and Video S3, Supporting Information). The total flow rate of the red and blue inner phases is fixed to obtain evenly spaced droplets at a constant printing speed. We program the distribution of the liquid inclusions in printing 3D objects. In the textured object in Figure 3d, the droplet spacing is precisely controlled in directions parallel and perpendicular to the printing path. This scheme of fine control on the droplet spacing could be modified for display applications. In Figure 3e, a quadrangle honeycomb architecture with monodisperse liquid inclusions is printed to demonstrate the ability of precise control over the inclusion spatial distribution. A printed block with three colors demonstrates the fast responses to the flow rate changing as well as to the mixing in red and blue inner phases (Figure 3f). In addition to printing structures with controlled inclusion distribution and composition with fast response time, this printing approach simplifies the preparation of composite inks and reduces the amount of waste material by direct dispersing inclusion materials at specified volumes.

To illustrate the potential of this multimaterial microfluidic 3D printing technique in producing functional composites, we fabricate a beam device which could undergo programmed shape changes when triggered by a thermal stimulus and print a self-healing polymer composite containing liquid inclusions filled with healing agents.

In the example of the beam device, we fabricate a two-layer beam with a resin top layer and a bottom layer consisting of liquid inclusions embedded in a resin matrix. As the thermal expansion properties of the liquid inclusions and the resin material are different, the top and bottom layers would differ in length due to the thermal mismatch strain and the beam bends. The effective bulk modulus \(\kappa_e\) and thermal expansion coefficient \(\alpha_e\) of the bottom composite layer are given by the micromechanics theory as

\[
\kappa_e = \frac{3\kappa_m\kappa_i + 4\mu_m[(1-c_i)\kappa_m + c_i\kappa_i]}{3[(1-c_i)\kappa_i + c_i\kappa_m] + 4\mu_m}
\]

(3)

\[
\alpha_e \kappa_i - \alpha_m \kappa_m = \frac{\alpha_i \kappa_i - [c_i \alpha_i \kappa_i + (1-c_i)\alpha_m \kappa_m]}{\kappa_i - \kappa_m}
\]

(4)

where \(\kappa, \mu,\) and \(\alpha\) represent the bulk modulus, shear modulus, and thermal expansion coefficient, respectively, and \(c_i\) is the volume fraction of the liquid inclusions. Subscripts “i” and “m” are used to identify quantities associated with the liquid inclusions and matrix, respectively.

Figure 3. 1D, 2D, and 3D structures printed by the multimaterial microfluidic 3D printing system. a) 1D spiral line with evenly spaced liquid inclusions. b) 2D layer with droplet inclusions in a pattern showing the letters PKU. c) 2D layer containing droplets of deliberately controlled compositions indicated by color gradient. d) Trapezoid block and e) 3D quadrangle honeycomb architecture with monodisperse liquid inclusions. f) Block with three steps, each consisting of inclusions filled with liquid of a specified color.
and matrix, respectively. Here we have assumed that the shear modulus of the liquid inclusions is negligible. The mechanical properties of the resin matrix are \( \mu_m = 0.36 \text{ GPa} \) and \( \kappa_m = 2.56 \text{ GPa} \), and \( \kappa_i = 2.20 \text{ GPa} \) for the water–glycerol inclusions. At \( c_i = 15\% \), the effective bulk modulus is determined as \( \kappa_e = 1.70 \text{ GPa} \) and the effective Young’s modulus is predicted as \( E_e = 2(1 - 2\nu)\kappa_e = 1.00 \text{ GPa} \) with Poisson ratio \( \nu = 0.40 \). As shown in Figure 4a, the theoretical prediction of \( E_e \) agrees well with the experimental result 0.96 GPa derived from the stress–strain curve in uniaxial tensile test. Thermal expansion tests indicate that the thermal expansion coefficients of the resin matrix and the composite (\( c_i = 15\% \)) vary with temperature nonmonotonically (Figure 4b). Negative thermal expansion coefficients near 70 °C for both the resin matrix and the composite might be due to the glass transition of the resin. With the knowledge of the effective Young’s modulus and thermal expansion coefficient of the composite, the thermal and mechanical behavior of the two-layer beam with a resin top layer and a bottom layer consisting of liquid inclusions embedded in a resin matrix can be determined and predicted. Here we perform experimental and theoretical analysis on such a two-layer beam fabricated using our microfluidic 3D printing system. The thickness of each layer in the printed two-layer beam is \( H = 150 \mu\text{m} \). Due to the different thermal expansion properties, the beam upon a temperature rise from \( T_0 = 26 \degree \text{C} \) to \( T_1 = 100 \degree \text{C} \) would bend into a circular arc of radius \( R \) as estimated by the beam theory on the thermal deflection of bimetal strip thermostat\(^{[37]} \) (Supporting Information) and is given by

\[
R = \frac{14 + E_m + E_e}{12H} \int_0^T (\alpha_e - \alpha_m) \, dT
\]

where \( \alpha_m \) and \( \alpha_e \) are the thermal expansion coefficients of the resin top layer and the composite bottom layer, respectively. Taking their values in Figure 4b, \( R/H \) in Equation (5) is estimated as 500 at 100 °C. Then we have \( R = 7.5 \text{ cm} \). The theoretical prediction (yellow dashed line in Figure 4c) agrees well with the experimental result. In Figure 4a–c, the volume fraction of liquid inclusions in the composite layer is \( c_i = 15\% \).

To investigate how the mechanical behaviors of the beam

![Figure 4](https://www.advancedscience.com)

**Figure 4.** Thermomechanical behaviors of printed composites made up of resin matrix with water–glycerol inclusions. a) The uniaxial tensile and b) thermal expansion tests of pure resin and resin-based composite with liquid inclusions at a volume fraction \( c_i = 15\% \). c) Thermal deflection of a printed two-layer beam (top layer, pure resin; bottom composite layer, resin containing liquid inclusions). The dashed line shows the theoretical prediction in Equation (5) with \( c_i = 15\% \). d) Bending of two-layer beams at different \( c_i \) from 0% to 25%. e) Thermally induced flower-shaped complex based on the two-layer beam structure. The beam could maintain at the flower-shaped configuration upon fast cooling to the room temperature. Scale bars, 1 cm.
depend on $c_i$ in the bottom layer, we fabricate two-layer beams of different $c_i$ and investigate their thermal deflection. As shown in Figure 4d, the two-layer beam exhibits larger deflection as $c_i$ increases which means that the effective thermal expansion coefficient of the composite bottom layer increases as $c_i$ increases. This is consistent with our theoretical prediction in Figure S6 (Supporting Information). As a featured demonstration, a smart structure with eight two-layer beams, which can switch reversibly between the flat and curved configurations upon thermal stimulus, is designed and fabricated (see Figure 4e). Taking advantage of the shape memory ability of the resin materials, the structure in the curved configuration could recover its original flat configuration upon slow cooling, and maintain the current configuration upon a fast cooling which can then undergo configuration recovery when reheated and slowly cooled. Such a smart structure has immediate potential applications in self-folding structures, structural composites, and soft actuator-based robots and devices.[38,39]

As a final demonstration on the versatility of this microfluidic-based 3D printing system, we fabricate a self-healing polymer composite consisting of liquid inclusions filled with two-part epoxy adhesive healing agents (Devcon, No. 14270) using the dual-nozzle printhead (Figure 5). Healing agents A (epoxy resin) and B (curing agent) are dyed red and blue, respectively, and are dispersed as liquid inclusions of the same size in the resin matrix. As the damage is inflicted on the composite material, the healing agents are released by the cracks crossing and fracturing them, and mixed at a volume ratio around 1:1. Within around 5 min local cracks are healed through mending reaction at the damage sites as demonstrated in Figure 5b,c, where a small crack of width around 58 $\mu$m and a relatively large crack of width around 500 $\mu$m are healed. Further tests indicate that the healed sample could carry an object of weight 200 g, showing the restoration of its mechanical strength (Figure 5c).

In conclusion, we have presented a multimaterial microfluidic 3D printing framework for fabricating textured composites with liquid inclusions of programmable distribution and compositions. Quantitative phase diagrams providing rational guide for the selection of desired printing parameters are determined. Using this microfluidic-based multimaterial printing framework, 1D, 2D, and 3D structures programmed with inclusion packing patterns and compositions are fabricated. We have also demonstrated the versatility of this 3D printing framework in the implications of producing layered composite beams and self-healing materials. The presented printing framework demonstrated using resin materials could be applied to a wide range of polymer materials and paves the way for the design and programming of composite material systems.

**Experimental Section**

**Microfluidic Chip Fabrication:** The microfluidic chips were designed using Solidworks (Dassault Systèmes) and fabricated with the aid of a stereolithography 3D printer (Form-2, Formlabs). The printed microfluidic chips were washed by isopropyl alcohol to remove uncured resin, then postcured for 1 h upon UV light (36 W), and eventually polished by fine sandpapers to achieve desired optical transparency.

**Printing System Setup:** The printing system (see Figure 1) is based on a commercial printer LULZBOT (Aleph Objects) with its original printhead replaced by the customized printhead fabricated using a commercial 3D printing system (Objet350 Connex3, Stratasys). Microfluidic chips and the print nozzle were assembled and added to the printhead. The cooperation of the ink pumping and the movement of the nozzle and the printing stage was achieved using the Arduino board (Zduino UNO) and Rambo board (Rambo rev 1.3 L). Commercial photosensitive resin (CLEAR PLG PCL04,

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**Figure 5.** Self-healing composites printed by the microfluidic 3D printing platform. a) Schematic of the self-healing process. The damage triggers the material rupture and release of the healing agents (epoxy adhesives A and B here) which then flow toward the damage sites and cause the local mending reaction to heal the crack. b) Self-healing of small cracks with a width around 58 $\mu$m. c) Self-healing of relatively large cracks with a width around 500 $\mu$m. The healed sample is able to carry a weight of 200 g.
Formlabs, density 0.98 kg cm⁻³) was used as the outer phase, and water–glycerol mixture (volume fraction 1:1, density 1.138 kg cm⁻³, viscosity 6.9 mPa s at 26 °C) dyed by colorant was used as the inner phase in a microfluidic flow. In the microfluidic chips, these two phases formed the ink consisting of water–glycerol inclusions embedded in the resin phase. The ink was extruded out of the blunt nozzle tip and was cured under a UV lamp (light wavelength 395 nm, output power 8 W) installed on the printhead.

**Contact Angle Measurement:** Pure resin droplets, each with a volume of 0.3 µL were deposited by a syringe on the glass substrate. The instrument Contact Angle Meter (Data Physics Corp.) was used to measure the resin contact angle. Data for each droplet were expressed as an average of five measurements.

**Line Printing:** In the printing of 1D lines (Figure 2a), the nozzle moved at different tip heights H from 170 to 500 µm and five different printing speeds v from 120 to 600 mm min⁻¹. The inner and outer phase flow rates Q₁ and Q₂ were fixed as 5 and 25 µL min⁻¹, respectively. The liquid inclusion sizes and spacing were analyzed by an image processing program ImageJ.

**Printing of 2D Carpet-Like Structure:** The nozzle height was fixed at 150 µm and the same printing speeds as those in Figure 2a were used. The spacing between the parallel straight nozzle motion paths from left to right in Figure 2b are 0.5, 0.7, 1.0, 1.5, 1.8, 2.0, 2.5, and 3.0 in the units of mm. The inner phase flow rates Q₁ were taken at 5, 5, 5, 5, 8, 6, 6, and 4, and the corresponding outer phase flow rates Q₂ were 5, 10, 15, 20, 18, 14, 20, and 6. All flow rates were in the units of µL min⁻¹. Inclination sizes and spacing were analyzed by ImageJ.

**Mechanical and Thermal Testing:** Tension tests of the printed samples were performed by the mechanical testing system MTI Instruments (MTESQuattro, ADMET Inc.). The thermal expansion coefficients were determined using the thermal testing device Dilatometer (DIL402C).

### Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

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### Conflict of Interest

The authors declare no conflict of interest.

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Supporting Information

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Figure S2 Rheological properties of the resin outer phase.

Figure S3 The relationship between the normalized droplet diameter \( D_{\text{drop}}/w_0 \) and the outer phase capillary number \( C_{\text{a}_o} \) at the T-junction in the generator of mixing droplets.

Figure S4 Structures of the printed line.

Figure S5 Inclusion positioning induced by the density difference between the inclusions and the fluidic matrix (resin here).

Figure S6 The effective Young’s modulus \( E_e \) and thermal expansion coefficient \( \alpha_e \).

Figure S7 Thermal deflection of the two-layer beam in a temperature rise.

Video S1 Printing a spiral structure.

Video S2 Printing the letters PKU (abbreviation of Peking University).

Video S3 Producing on-demand mixing droplets.

For the two-layer beam upon thermal deflection in Figure 4 in the main text, the top layer is composed of pure resin and the composite bottom layer consists of liquid inclusions suspended in the resin matrix as illustrated in Figure S7. Each layer has a thickness of \( H \). The thermal expansion coefficient of the bottom layer is larger than that of the top layer as shown in Figure 4b. The initial straight beam has a length \( L_0 \) at the room temperature \( T_0 \). An increase in temperature to \( T_1 \) induces a relative length increase \( \Delta L_m = \int_{T_0}^{T_1} \alpha_m dT \) in the bottom and a relative length increase \( \Delta L_e = \int_{T_0}^{T_1} \alpha_e dT \) in the top layer, and the beam deforms into a circular arc of an arc angle \( \theta \) and a radius of curvature \( R \). Here \( \alpha_m \) and \( \alpha_e \) are the thermal expansion coefficients of the resin top layer and the composite bottom layer, respectively. Based on the beam theory on the thermal deflection of a bi-metal strip termostat, the normalized relative length increase of the interface between the top and bottom layers can be given as

\[
\varepsilon_{\text{it}} = \int_{T_0}^{T_1} \alpha_m dT + \frac{H}{12R} \frac{E_m + E_e}{E_m} + \frac{H}{2R} \quad \text{or} \quad \varepsilon_{\text{ib}} = \int_{T_0}^{T_1} \alpha_e dT - \frac{H}{12R} \frac{E_m + E_e}{E_e} - \frac{H}{2R}.
\]

Then we have \( \varepsilon_{\text{it}} = \varepsilon_{\text{ib}} \) or

\[
\frac{R}{H} = \frac{14 + \frac{E_m + E_e}{E_m}}{12 \int_{T_0}^{T_1} (\alpha_e - \alpha_m) dT}.
\]

Based on the experimental data on \( E_m \) and \( E_e \) in Figure 4a and \( \alpha_m \) and \( \alpha_e \) in Figure 4b, \( R/H \) is then estimated as 500. As \( H=150 \mu m \), we have \( R=7.5 \) cm which agrees well with the experimental result in Figure 4c. The thickness variation of each layer during the thermal deformation has been ignored.
**Figure S1.** T-junction droplet generator. The upper left inset shows the assembled chip in the printing platform and the upper right inset indicates the on-demand mixing droplet formation.

**Figure S2.** Rheological properties of the resin outer phase.

**Figure S3.** The relationship between the normalized droplet diameter $D_{\text{drop}}/w_o$ and the outer phase capillary number $C_{a_o}$ at the T-junction in the generator of mixing droplets. Symbols represent the experimental data and solid lines show the scaling law $D_{\text{drop}}/w_o \sim C_{a_o}^{-1}$. 
Figure S4. Structures of the printed line. (a) The line width \( w \) as a function of \( \sqrt{Q/v} \) with \( Q \) as the total flow rate and \( v \) as the printing speed. Upper inset: schematic of a printed line of width \( w \), height \( h \), and contact angle \( \theta \). Lower inset: contact angle of the resin on the glass substrate at room temperature. (b) The linear relationship between the normalized droplet spacing \( \Delta x/D_{\text{nozzle}} \) and the dimensionless parameter \( V_{\text{droplet}} v/(D_{\text{nozzle}} Q) \) with \( V_{\text{droplet}} = \pi D_{\text{drop}}^3 / 6 \). Dots in (a) and (b) are the experimental results and dashed lines represent the theoretical prediction.

Figure S5. Inclusion positioning induced by the density difference between the inclusions and the fluidic matrix (resin here). (a) Illustration of the positioning of the droplet inclusions. (b) Droplets are stable and remain intact in a helical structure formed by pulling the 1D spiral line in Figure 3a in the main text. (c) In the case of the bubble inclusions, the bubbles of a much smaller density than the matrix first rise and then break at the resin surface.
**Figure S6.** The effective Young’s modulus $E_e$ (a) and thermal expansion coefficient $\alpha_e$ at 100 °C (b) of the two-layer beam as functions of the droplet volume fraction.

**Figure S7.** Thermal deflection of the two-layer beam in a temperature rise. The initial straight beam has a length $L_0$ at the room temperature $T_0$. An increase in temperature to $T_1$ induces a larger length increase $\Delta L_e$ in the bottom layer than the length increase $\Delta L_m$ in the top layer, and the beam becomes curved with a radius of curvature $R$. 