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High-Speed Vapor Transport Deposition of Perovskite Thin Films

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Supporting Information

ABSTRACT: Intensive research of hybrid metal-halide perovskite materials for use as photoactive materials has resulted in an unmatched increase in the power conversion efficiency of perovskite photovoltaics (PVs) over the last couple of years. Now that lab-fabricated perovskite devices rival the efficiency of silicon PVs, the next challenge of scalable mass manufacturing of large perovskite PV panels remains to be solved. For that purpose, it is still unclear which manufacturing method will provide the lowest processing cost and highest quality solar cells. Vapor deposition has been proven to work well for perovskites as a controllable and repeatable thin-film deposition technique but with processing speeds currently too slow to adequately lower the production costs. Addressing this challenge, in the present work, we demonstrate a high-speed vapor transport processing technique in a custom-built reactor that produces high-quality perovskite films with unprecedented deposition speed exceeding 1 nm/s, over 10× faster than previous vapor deposition demonstrations. We show that the semiconducting perovskite films produced with this method have excellent crystallinity and optoelectronic properties with 10 ns charge carrier lifetime, enabling us to fabricate the first photovoltaic devices made by perovskite vapor transport deposition. Our experiments are guided by computational fluid dynamics simulations that also predict that this technique could lead to deposition rates on the order of micrometers per second. This, in turn, could enable cost-effective scalable manufacturing of the perovskite-based solar technologies.

KEYWORDS: solar cells, perovskite, thin-film, vapor deposition, manufacturing, fluid dynamics

INTRODUCTION

The latest report showing hybrid metal-halide perovskite solar cells with power conversion efficiency (PCE) exceeding 23% for longer than 1000 h−1 is a remarkable demonstration of one of the presently most promising emerging energy technologies. In this record demonstration, as well as in other recent reports of the highest-PCE perovskite solar cells, the photoactive layers of the devices were fabricated via spin-coating method, which, unfortunately, is not scalable, as would be needed for their mass production. Solution processing brings a number of challenges inherent in the solvent drying kinetics, which make it difficult to independently control film formation and crystallization to achieve the desired film quality, especially over large areas. The nonuniformity of solution-coated films can introduce macroscopic defects that can act as shunt pathways, nonactive areas, and avenues for moisture and metal ingress, compromising device stability. Solution processing also introduces fabrication challenges for more advanced stacked multilayer perovskite solar cells as solvents used in one layer need to be compatible with deposition on top of all of the underlying layers. Above all, standard solvents used in perovskite film processing are often toxic, resulting in hazardous solvent waste on a production line and making it to either replace them or eliminate them, as we do in the present work where we demonstrate solvent-free vapor deposition.

State-of-the-art vapor-deposited perovskite solar cells, fabricated via vacuum thermal evaporation (VTE), have exceeded 20% efficiency. However, this record demonstration relied on the photoactive layer deposition at low operating pressures (10−6 Torr), requiring ultra-high-vacuum equipment and long pump-down times that could be limiting manufacturing throughput. The speed of production influences heavily the manufacturing cost of the solar module. As the price for currently sold photovoltaics (PV) modules is well below 1 USD/Wp, the deposition speed needs to be at values higher than currently demonstrated VTE rates of Å/s. Additionally, the relatively low material utilization efficiency of <50% increases both material cost per watt and chamber cleaning requirements. Since the evaporation rate varies exponentially
with and can only be controlled by the source temperature, \(^{13}\) precise optimization of multiple co-deposited precursors, as needed for perovskite film deposition, would be difficult to control, and the needed high deposition rates would be limited by the relatively low decomposition temperature of the material.

Many of these limitations and drawbacks could be overcome with another vapor-based process known as vapor transport deposition (VTD), which operates under minimal low-pressure conditions. VTD relies on a transport carrier gas to carry sublimated material vapors from the source to the target substrate, not requiring a chemical reaction of the materials in the vapor phase. VTD has been used commercially to deposit high-quality CdTe films for solar applications at high rates (≥500 nm/s).\(^ {14,15}\) In the field of organic semiconducting materials, a similar process named “organic vapor phase deposition” (OVPD) has been developed, in which organic films are deposited from evaporated organic small molecules delivered by a gas stream through a heated apparatus onto a substrate.\(^ {16}\) It has been demonstrated that this process can be operated at five orders of magnitude higher pressure as compared to thermal evaporation, which is not only saving costs but allows to independently control the deposition rate from temperature by controlling the gas flow rate and pressure. This inherently leads to better control over the morphology of the grown thin films, hence allowing for tuning of material properties.\(^ {17−19}\)

In the field of perovskite solar cells, chemical vapor deposition processes in tube furnaces have been used to convert solution-deposited or thermally evaporated lead halide (PbI\(_2\) or PbCl\(_2\)) films into perovskites by exposing them to vapor-transported methylammonium iodide (MAI) or formamidinium iodide.\(^ {20−22}\) A similar result was reported by combined chemical vapor deposition of MAI and PbI\(_2\) in the same vacuum tube.\(^ {23}\) The demonstrated results of high-quality perovskite films over large areas prove that vapor-based scaling methods represent a promising route toward forming scalable perovskite solar cells. However, these first demonstrations of precursor depositions and conversion to perovskite films in tube furnaces were unable to exceed the deposition rates of conventional thermal evaporation-based deposition demonstrations, which are inferior to solution-based methods.

Here, we report on the development of a novel method for depositing perovskite thin films with unprecedented high throughput. Our solvent-free process is based on vapor transport deposition with customized adaptations for specific perovskite requirements. We designed dedicated material sublimation sources for the precursor materials that are transported and guided in the vapor state by the individually controlled carrier gas flow toward a temperature-controlled substrate. The fully vapor-deposited perovskite thin films are formed by combining multiple precursors in an alloyed polycrystalline film. We show that perovskite VTD requires pressures not lower than 10 Torr, achievable with simple pumping and chamber equipment. This reduces capital expenditures substantially compared to high-vacuum thermal evaporation. At these pressures, precursor molecules are randomly distributed in the carrier gas, combined with precisely controlled flow, this allows for highly uniform film deposition with high material utilization. VTD decouples evaporation and condensation, allowing orders-of-magnitude higher deposition rates without decomposing the materials.

Accurate compositional control can enable precise stoichiometric deposition of multiple compounds, which, in turn, permits the fabrication of the multicomponent perovskite films. Crystallization conditions can be fine-tuned by varying the substrate temperature and deposition pressure, allowing control of the film morphology and defect densities, a key requirement for making efficient and stable devices, hence,
RESULTS AND DISCUSSION

We developed a multisource VTD system that integrates temperature and material control with the gas flow (Figure 1a), using nitrogen as the carrier gas. The nitrogen gas is first heated by passing it through heated tubes that are sufficiently long to ensure uniform gas temperature, and then it flows through the equally heated vaporization zones containing the precursor materials, where they get sublimated into the vapor phase and then carried by the transport gas toward a temperature-controlled substrate. The sublimation itself consumes energy, which could lead to a cooling of the system. However, the amount of material being vaporized is low enough that we can assume no influence on the gas flow. The magnified cross-section of the diagram in Figure 1b illustrates the inner components of the VTD system. The custom-made capsules within the source tubes contain the source material, which is capped by two silicon carbide (SiC) foam filters. The capsules with foam filters can be easily extracted and cleaned for optimal uncontaminated operation. The lower filter is used to prevent the solid materials from escaping the source tube, whereas the upper filter ensures uniform gas flow and temperature distribution of the transport gas within the capsule. A temperature-controlled liquid is circulated through the substrate stage to allow for precise control of the substrate temperature, which is critical in the control of the crystallinity of deposited films. The deposition chamber operating pressure and gas flow are symmetrical. Figure 2a shows the temperature map in the top half and the pressure field in the bottom half. We notice from the plots that the temperature varies substantially throughout the geometry. The high flow rate influences the temperature distribution at the inlet of the PbI₂ source tube, which is wrapped with heating tape and well controlled. However, the source cell shows perfectly uniform heating at the constant maximum temperature of 450 °C, due to the slightly slowed downflow through the SiC filters. The pressure is almost constant at the controlled value of 10 Torr throughout the geometry and only undergoes a slight increase through the filters at steady state. Figure 2b shows the concentration map in the top half and the velocity map with field arrows in the bottom half.

We additionally calculated the deposition rate on the chamber walls and on the substrate stage and plotted it as a color contour edge within the concentration map. We notice that the concentration map follows the shape of the velocity field. Hence, the distribution of the highly diluted PbI₂ vapor within the N₂ carrier gas is mostly driven by the directional laminar flow. Only at the boundary layers, it is driven by the concentration gradient, which is influenced by the temperature gradient. We find that the deposition rate (shown as color-mapped boundaries in Figure 2b) at the center of the substrate comes to roughly 1.5 nm/s and increases slightly at the edges, similar to the OVPD simulation results shown by others. The equivalent modeling results of the MAI deposition are plotted in Figure S1. Due to the high vapor pressure of MAI, the deposition rate for a co-deposition is limited by the PbI₂ deposition and potentially the interdiffusion and crystal growth kinetics that have not yet been considered. When analyzing the
deposition rate, we encounter very high parasitic deposition of the material within the part of the source tube after the source cell. The current design does not allow for heating the end of the center tube to higher temperatures. For an improved design, heating the reactor walls to the substrate stage should be considered to maximize material utilization and to accomplish higher deposition rates. In such a case, the substrate stage would have to be substantially cooled to keep the surface temperature low enough for efficient deposition. The results of this simulation should be used predominantly for qualitative trend analysis instead of exact quantitative predictions, as they were computed in simplified 2D space due to computation time limitations and lacking an exact sticking coefficient study for precise deposition rate analysis. Both are beyond the scope of this work.

Guided by the simulation results, we sequentially deposited perovskite films by first vapor-transporting PbI_2 onto the substrate and then subsequently exposing it to a transported vapor stream of MAI to form MAPI perovskite films. This deposition sequence allows us to adjust the right stoichiometric compositional combination of the materials by changing only the deposition times. To optimize the process and compare deposition rates with simulation results, we started with different deposition time durations for the two materials. After preheating the source tubes for at least 5 min (PbI_2) or 2 min (MAI) to allow for steady-state sublimation, we opened the nitrogen carrier gas valves for different deposition times. We show the appearance of the increasingly thicker PbI_2 films in the inset photographs of Figure 3a and plot their thicknesses, determined by a profilometer, as a function of the increasing deposition durations. The resulting linear growth behavior corresponds to a growth rate of 1.3 nm/s, which is about 10 times faster than demonstrations of thermally evaporated PbI_2 films. We notice a difference between the increasing thickness of the deposited films toward the edge and the deposition rate profiles of the simulation. As the deposition rate is highly dependent on the boundary layer (velocity profile at the deposition surface), it is likely that our simplified 2D simulation does not lead to correct results of the velocity simulation between the round inner chamber walls and the quadratic substrate stage (visible in the corners of the substrates with VT-deposited films). Furthermore, it is plausible that the temperature of our substrates is not uniform, as assumed in the simulation, but hotter at the edges due to reduced heat distribution within the substrate stage. This would lead to a significantly reduced deposition rate, as the sticking coefficient has a strong dependence on temperature. We decided to use a PbI_2 film thickness of ca. 200 nm, reached after 3 min of VTD, to convert to MAPI films. As discussed earlier, the PbI_2 deposition rate is the limiting process as the vapor pressure of MAI is substantially higher, which corresponds to a much faster deposition speed. The increased duration of MAI VTD is leading to films with initial PbI_2 excess and then gradually changing to more stoichiometrically balanced compositions and then to compositions with an excess of MAI. This coarse optimization process is displayed as a color change in the inset photographs of Figure 3a. We examined the morphology of the optimal VTD PbI_2 film and converted MAPI film by 50× magnified optical microscopy and scanning electron microscope (SEM) imaging of the surfaces. The resulting images of PbI_2 (b, d) and MAPI (c, e) are shown in Figure 3b–e. We find that the PbI_2 morphology appears to be porous with ca. 100 nm large islands, which could help to allow for enhanced MAPI conversion due to larger available surface area. Consequently, the pin-hole free morphology of MAPI appears to have a good quality perovskite crystal grain size of roughly 200 nm. Hence, these films should allow for well-functioning absorber layers in perovskite solar cell devices with less trap probability within the grain boundaries.

 Sequentially depositing the precursor materials is not regarded as the fastest and, therefore, most scalable technique for the vapor-based perovskite fabrication. This is due to the conversion of PbI_2 to MAPI upon MAI exposure requiring an initial diffusion of the organic material to then convert to the perovskite crystal structure. Thick PbI_2 (>100 nm) films, required to make thick enough absorber layers to maximize charge carrier generation, require a longer time for the intercalation of MAI and conversion to perovskite. One solution to this challenge is to deposit an alternating sequence of thin pairs of films that add up in thickness to the original thick film pair. Specifically, we deposited pairs of PbI and MAI three times in a row. To compare the optical properties of the resulting films, we measured the UV-vis absorbance spectra of bare PbI films that were deposited via VTD as well as a sequentially deposited layer of PbI_2 for 3 min and then vapor transport (VT) depositing MAI for 1 min and repeating the former procedure three times. Results shown in Figure 4a
demonstrate that the alternating sequential VTD technique leads to perovskite films that absorb significantly more light. The absorbance onset suggests a band-gap edge at a wavelength of 780 nm (corresponding to 1.59 eV), which is in accordance with the typical properties of MAPI films. The VTD PbI₂ film shows a band-gap edge at 520 nm (corresponding to 2.39 eV), which is also in good agreement with other reports. To analyze the detailed compositional changes in the different MAPI films, we measured the X-ray diffraction (XRD) patterns, as shown in Figure 4b. The three times alternated sequentially deposited perovskite films are more crystalline, as evidenced by the nearly complete diminishment of the lead iodide peak, which has a typical appearance at around 12°-13° and the dominance of the MAPI perovskite peak at 14°. A trial of four times alternated sequentially deposited perovskite films resulted in substantial excess of MAI on the surface, which lead to rapidly degrading films that could not be measured quickly enough. On an additional sample, we reduced the MAI deposition time slightly to see the effect in the XRD spectrum and find that the conversion to the perovskite phase is reduced, as determined by the decreased MAPI peak at 14° and the reappearing of the PbI₂ peak. Reassuring us of having reached a close to optimal condition with the initial protocol of three times alternated sequential deposition.

Figure 4c plots the PL decays of a bare VTD MAPI film together with a fitted stretched exponential decay. We find that the average monomolecular carrier lifetime of our VTD CH₃NH₃PbI₃ perovskite film, determined with the exponential fit, equals 10.3 ns. Although this is shorter than some results taken from films made via solution processing or special passivation, it appears to match with published results from thermally evaporated films without passivation that leads to well-working solar cell devices.

Above results indicate that VTD MAPI films have good optoelectronic properties and good morphology, which could lead to well-performing MAPI perovskite solar cells. To fabricate solar cells with VTD MAPI films, we start with fluorinated tin oxide (FTO)-coated glass on top of which a TiO₂ compact electron transporting layer is spin-coated and high-temperature-annealed. The VTD MAPI layer is then deposited and capped with a spin-coated spiro-OMeTAD hole-transporting layer and a thermally evaporated gold top contact. The device stack is shown in the inset cross-sectional SEM image of Figure 4d, with the J−V characterization under simulated AM 1.5 sunlight plotted in Figure 4d. The functioning device has a PCE of 6.9%, coming from a short-circuit current (J_SC) of 14.2 mA/cm², and a notably good open-circuit voltage (V_OC) of 1.01 V, indicating low recombination losses due to defect states or surface recombination. The fill factor (FF) of 0.48 is significantly lower than current high-performance results, and the shape of the curve indicates that there are substantial parasitic series resistances. Together with the relatively low J_SC, this indicates very inefficient charge extraction at one of the interfaces, which could be caused by an energy barrier from unoptimized perovskite compositions that still contain an excess of MAI or PbI₂ and need to be improved by further optimizing the growth conditions. Additionally, the low J_SC and FF could be an indication of a poor charge collection due to the
Figure 5. Simulated deposition rate as a function of source temperature and flow rate at chamber pressures of 0.1 Torr (a), 1 Torr (b), and 10 Torr (c).

nonoptimized thickness and material quality of the charge-transport layers. The device also displays a severe hysteresis effect when comparing the $J_V$ curves from forward ($J_{SC}$ to $V_{OC}$) and reverse ($V_{OC}$ to $J_{SC}$) scan directions, as shown in Figure S3. The hysteresis is most likely an effect of the device architecture, as devices with MAPI as the absorber material and Li-doped spiro-OMeTAD as the hole-transporting material as well as TiO$_2$ as the electron transporting material have been demonstrated to display this effect. We did not carry out any device stability studies, since devices that use this combination of materials are well known to be prone to quick degradation under ambient conditions.

The scalability of the perovskite film fabrication with this novel VTD technique is the core focus of the present work. In the present implementation, the perovskite film-growth rate is dependent on the deposition rate of PbI$_2$ films. Therefore, to explore the deposition rate limits, we changed the processing conditions within our computational model to allow for optimized elevated source tube heating and chamber-wall heating to minimize the parasitic chamber-wall deposition during the film-growth process and maximize PbI$_2$ utilization. We show the effect that the processing parameters have on the deposition rate, by plotting the latter for varied source temperature and carrier gas flow rate at three different chamber pressures of 0.1 Torr (Figure 5a), 1 Torr (Figure 5b), and 10 Torr (Figure 5c). To check if the simulation can be appropriately carried out in the continuum flow regime even at lower pressures, we calculated the maximum Knudsen effect that the processing parameters have on the deposition rate, easily explained by the exponential dependence of sublimation on temperature. The limit would be reached at the decomposition temperature of PbI$_2$.

It still needs to be proven that this technique can produce devices with equally high performances, as can be achieved as with other deposition methods. Additionally, some engineering challenges could emerge due to the need for precise process control systems, moderately high-temperature resistant components, and the need for a cooled substrate stage. However, the first experimental demonstrations and the results of the simulations are promising and open the doors for a remarkably fast deposition method for the active layers of perovskite solar cells.

### EXPERIMENTAL SECTION

The glass substrates have been purchased from Pilkington Inc. with a coated fluorinated tin oxide (FTO) layer 15 Ω/sq. FTO was partially removed from the substrate via etching with zinc powder and 2 M HCl to create nonconductive stripes for later needed counter-electrode fabrication. The etched substrates were rinsed in deionized (DI) water and cleaned rigorously by brushing the Helmanex solution with a toothbrush before rinsing again in a stream of DI water. The substrates were then subsequently rinsed in acetone and 2-propanol and then etched for 10 min in O$_2$ plasma. The n-type TiO$_2$ compact layer was formed by spin-coating a solution of 0.71 g titanium isopropoxide and 0.07 g of 2 M HCl in 8 ml of ethanol with 2000 rpm for 45 s and then annealed at 500 °C for 45 min.

The perovskite thin films were deposited with the custom-built tool, as demonstrated in the main text. The procedure started with the opening of the vacuum-tight system and the loading of the source materials. To keep the system clean and to support the reproducibility of the experiments, the material capsules including the metal parts and the SiC filter disks are cleaned thoroughly by immersion into dimethylsulfoxide (DMSO) and sonication for 5 min and then immersing and rinsing in isopropanol before drying with a stream of nitrogen and placement onto a 150 °C hotplate for 10 min. After assembly of the capsules, roughly 4 g of lead iodide (PbI$_2$) powder was loaded into one capsule and ca. 4 g of methylammonium iodide (MAI) powder was loaded into a second capsule. The source tubes of the VTD tool were cleaned by wiping the inside with DMSO-soaked cleanroom wipes that were pushed through the tube by a metal rod. The lead iodide capsule was then placed and pushed from the N$_2$ supply side into the middle source tube until it would not move further due to a tube narrowing at the very end. The MAI capsule was placed into the left source tube in the same way. The subsequent loading of the previously prepared substrate was performed by using a Kapton tape, which covers only a small area on the side of the
substrate, to fasten it to the temperature-controlled metal stage. The tool was then closed up to allow for a roughing vacuum. The vacuum pump was turned on, and the system pressure reaches values close to $1 \times 10^{-3}$ Torr after a couple of minutes. The vacuum pump is kept clean by using a cold trap that needs to be refilled with liquid nitrogen before every use. To keep the pressure of the system constant, the side $N_2$ inlet valve was opened, and the flow adjusted so that the chamber reaches a pressure of 10 Torr. Then, the circulation bath heater was turned on to 125 °C, and after a couple of minutes, the temperature-controlled substrate stage reached a measured temperature of ca. 100 °C. Without the use of the circulation bath heater, the substrate stage reaches a maximum temperature of 70 °C. Therefore, we actively heated the substrate during deposition. Subsequently, the middle PbI$_2$ source tube is heated to 450 °C, which is controlled by adjusting the power source of the resistive heating band and probing the temperature at the outside of the tube. Once the target temperature is reached, we wait 5 min for uniform, steady-state temperature distribution, and sublimation of the material before we turn on the $N_2$ source flow and simultaneously turning off the side inlet. The deposition flow is adjusted by sensing the pressure to equal exactly 10 Torr during deposition. The film thickness is controlled by timing the deposition time. For best results, a deposition time of $3 \times 1$ min for the PbI$_2$ deposition was used. After the PbI$_2$ deposition is finished, we turn off the $N_2$ source flow and simultaneously turning on the side $N_2$ flow to keep the pressure constant. At the same time, we turn off the power source for the heating band to decrease the temperature of the middle source tube. Then, we repeat the procedure in a similar manner with the MAI source tube. There we set the target source temperature to 210 °C and only wait 2 min to stabilize before starting with the deposition, which is also performed with adjusted $N_2$ flow to keep 10 Torr system pressure. The best results are achieved with a deposition time of $3 \times 30$ s. After turning off the source flow while turning on the side port flow and letting the temperature decrease, we wait for 15 min before turning off the vacuum pump and opening the system to remove the substrate.

For the fabrication of solar cell devices, the hole-transporting layer was deposited on top of the perovskite films by spin-coating a solution with 8.5 wt % spiro-OMeTAD in chlorobenzene and 30 mol % lithium bis(trifluoromethane-sulfonyl)imide and 80 mol % 4-tert-butylpyridine as additives, at 2000 rpm for 60 s in the dry air atmosphere. The sample was dried overnight in a light-sealed desiccator. Au (100 nm) was thermally evaporated on top of spiro-OMeTAD as the top electrode. The pressure was ca. $10^{-3}$ Torr and the evaporation rate was close to 0.1 nm/s.

For measuring the performance of the solar cells, simulated AM 1.5 sunlight was generated with a Newport solar simulator containing a Spectra Physics light source by Oriel Instruments. It was calibrated to give simulated AM 1.5, of 100 mW/cm$^2$ equivalent irradiance, using a calibrated silicon reference diode with a filter. The current–voltage curves were recorded with an Oscilla source measure unit T2002D. The solar cells were masked with a metal aperture defining the active area (0.2 cm$^2$) of the solar cells.

For the thickness measurements, we used a Bruker DXT Stylus Profilometer. The SEM images were taken with a Zeiss FESEM Ultra55 and the optical microscope images with a Nikon LV100 and a 50X Nikon (NA/WD: 0.80/1.0 mm) objective. We used a Rigaku SmartLab X-Ray Diffractometer to measure the X-ray diffraction patterns and a Cary 5000 UV–Vis–NIR Spectrophotometer for the UV–vis spectra. For the time-resolved photoluminescence measurements, a 405 nm Picoquant laser was used, with 100 s of 1 ps pulse width and a repetition rate of up to 80 MHz.

Additional details about the computational methods applied for the simulations is provided in the Supporting Information.

**CONCLUSIONS**

We present a novel vapor transport deposition method for perovskite thin films enabled by a custom-made multiple-source tool. Given the multicomponent nature of perovskite thin films, we found that a precise process control is needed for deposition of high-quality perovskite films. The choice of our processing parameters was guided by computational fluid dynamic simulations, which allowed for a determination of the effects of carrier gas flow and temperature, on the chamber pressure and vapor concentration, further allowing us to predict the perovskite film deposition profile on the substrate and chamber walls. In our experiments, we arrived at very high film deposition rates, surpassing 1 nm/s for PbI$_2$ and even faster for MAI. We demonstrated that the alternating sequential deposition can lead to perovskite thin films with high-quality morphology and crystallinity, similar to other deposition techniques that yield state-of-the-art films. The first solar cells containing perovskite photoactive films deposited with this novel technique already show a PCE of 6.9%. Finally, our computational simulations indicate that deposition rates on the order of micrometers per second should be possible by further optimizing this VTD process, which would benefit future scalable production of perovskite solar cells.

**ASSOCIATED CONTENT**

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.9b07651.

Computational methods, process simulation; domain governing equations used in the CFD simulations; calculation of the Knudsen number; vacuum pump displacement curve at varying pressure; characteristic JV-curve of VT deposited perovskite solar cell device (PDF)

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**Notes**

The authors declare no competing financial interest.

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