Marginally trapped surfaces in $\mathbb{L}^4$ and
an extended Weierstrass-Bryant representation

Juan A. Aledo$^a$, José A. Gálvez$^b$ and Pablo Mira$^c$

$^a$ Departamento de Matemáticas, Universidad de Castilla la Mancha. Escuela Politécnica
Superior de Albacete, E-02071 Albacete, Spain
e-mail: juanangel.aledo@uclm.es

$^b$ Departamento de Geometría y Topología, Universidad de Granada, E-18071 Granada,
Spain.
e-mail: jagalvez@ugr.es

$^c$ Departamento de Matemática Aplicada y Estadística, Universidad Politécnica de
Cartagena, E-30203 Cartagena, Murcia, Spain.
e-mail: pablo.mira@upct.es

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Abstract

We give a conformal representation in terms of meromorphic data for a certain
class of spacelike surfaces in $\mathbb{L}^4$ whose mean curvature vector verifies $\langle H, H \rangle = 0$.
This representation extends simultaneously the Weierstrass representation for
minimal surfaces in $\mathbb{R}^3$ and for maximal surfaces in $\mathbb{L}^3$, and the Bryant repre-
sentation for mean curvature one surfaces in the hyperbolic 3-space and in the de
Sitter 3-space.

1 Introduction

In 1987 R.L. Bryant [Bry] described a conformal representation for the surfaces with
constant mean curvature $H = r$ in the hyperbolic 3-space $\mathbb{H}^3(-r^2)$ of constant curvature
$-r^2$, similar in spirit to the classical Weierstrass representation for minimal surfaces in
$\mathbb{R}^3$. Having Bryant’s work as a basis, the theory of CMC-$r$ surfaces in $\mathbb{H}^3(-r^2)$ has

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experimented in the last fifteen years a great development, largely influenced by the global results of minimal surface theory.

Additionally, Umehara and Yamada [UmYa] showed that there is an isometric deformation process via which CMC-\(r\) surfaces in \(\mathbb{H}^3(-r^2)\) converge analytically to minimal surfaces in \(\mathbb{R}^3\) as \(r \to 0\). This suggested the possibility of unifying both conformal representations into a more general one in a natural way.

The first such extended representation was obtained in [KTUY] for a class of surfaces with holomorphic right Gauss map in certain Lie groups containing those of the form \(\text{SL}(n,\mathbb{C})/\text{SU}(n)\). An alternative unified representation was derived in [HMN] in terms of Möbius geometry of surfaces.

The present paper provides a new conformal representation generalizing simultaneously those of Weierstrass and Bryant. To do so, we consider the Minkowski spacetime \(\mathbb{L}^4\) as the ambient space, and we view \(\mathbb{R}^3\) and \(\mathbb{H}^3(-r^2)\) as hyperquadrics of \(\mathbb{L}^4\) in the usual way. With this, we will describe a complex representation for a certain class of spacelike surfaces in \(\mathbb{L}^4\), which we will call surfaces of Bryant type in \(\mathbb{L}^4\), that includes the minimal surfaces in \(\mathbb{R}^3\) and the CMC-\(r\) surfaces in \(\mathbb{H}^3(-r^2)\). Indeed, the surfaces of Bryant type in \(\mathbb{L}^4\) that lie in \(\mathbb{R}^3 \subset \mathbb{L}^4\) (resp. \(\mathbb{H}^3(-r^2) \subset \mathbb{L}^4\)) are exactly the minimal surfaces of \(\mathbb{R}^3\) (resp. the CMC-\(r\) surfaces of \(\mathbb{H}^3(-r^2)\)).

The main geometric property of these Bryant-type surfaces is that their mean curvature vector \(H\) verifies \(\langle H, H \rangle = 0\), where \(\langle , \rangle\) is the Lorentzian product of \(\mathbb{L}^4\). The spacelike surfaces defined by this condition are well known in General Relativity, where they are called marginally trapped surfaces, and represent useful objects in the theory of singularities in spacetimes (see [HaEl]). From our viewpoint, the isotropy condition \(\langle H, H \rangle = 0\) implies that a certain Gauss map of the surface in \(\mathbb{L}^4\) is conformal, what generalizes the well known fact that both minimal surfaces in \(\mathbb{R}^3\) and CMC-\(r\) surfaces in \(\mathbb{H}^3(-r^2)\) have conformal Gauss maps.

There are some points of special interest in the present unified conformal representation. First, it does not only generalize the representation formulae in the theories of minimal surfaces in \(\mathbb{R}^3\) and CMC-\(r\) surfaces in \(\mathbb{H}^3(-r^2)\). It also includes the conformal representations of their Lorentzian counterparts, namely, the theories of maximal surfaces in \(\mathbb{L}^3\) [Kob] and of spacelike CMC-\(r\) surfaces in the de Sitter 3-space \(S^3_1(r^2)\) [AiAk], when we view \(\mathbb{L}^3\) and \(S^3_1(r^2)\) as hyperquadrics of \(\mathbb{L}^4\) in the usual way. In addition, with the present complex representation the Umehara-Yamada perturbation process is to some extent simplified, as it is viewed in the fixed ambient space \(\mathbb{L}^4\). Finally, the conformal representation can be used to construct many complete surfaces of Bryant type in \(\mathbb{L}^4\) which do not belong to any of the previous families, but that still have physical interest as they are marginally trapped surfaces in \(\mathbb{L}^4\).

The paper is organized as follows. In Section 2 we analyze the structure equations of a spacelike surface \(\psi : \Sigma \to \mathbb{L}^4\), and prove that a natural hyperbolic Gauss map \(G : \Sigma \to \mathbb{C} \cup \{\infty\}\) on the surface is conformal if \(\psi\) is a marginally trapped surface. We also show that if the normal bundle of \(\psi\) is flat, then a certain Hopf differential on the surface is holomorphic.

Section 3 describes the basic result of the present work: a conformal representation
for the surfaces of Bryant type in $\mathbb{L}^4$. Here we define a *surface of Bryant type* in $\mathbb{L}^4$ as a marginally trapped surface with flat normal bundle that is locally isometric to a minimal surface in $\mathbb{R}^3$ or to a maximal surface in $\mathbb{L}^3$.

In Section 4 we will show that the meromorphic representation we have obtained generalizes the Weierstrass representation of the minimal (resp. maximal) surfaces in $\mathbb{R}^3$ (resp. $\mathbb{L}^3$), and the Bryant representation of the CMC-$r$ surfaces in $\mathbb{H}^3(-r^2)$ and $S^3_1(r^2)$. We will also indicate how the Umehara-Yamada deformation is described in our context, and we will construct new examples of complete surfaces of Bryant type in $\mathbb{L}^4$ that do not belong to any of the previous families.

Finally, in Section 5 we will classify the complete surfaces of Bryant type in $\mathbb{L}^4$ with non-negative curvature, as well as the complete simply-connected surfaces of Bryant type in $\mathbb{L}^4$ with finite total curvature. The paper ends up with an appendix containing some auxiliary results.

## 2 Marginally trapped surfaces

Let $\mathbb{L}^4$ denote the 4-dimensional Lorentz-Minkowski space, that is, the real vector space $\mathbb{R}^4$ endowed with the Lorentzian metric

$$\langle , \rangle = -dx_0^2 + dx_1^2 + dx_2^2 + dx_3^2,$$

in canonical coordinates. We shall identify $\mathbb{L}^4$ with the space of 2 by 2 Hermitian matrices in the usual way,

$$(x_0, x_1, x_2, x_3) \in \mathbb{L}^4 \longleftrightarrow \begin{pmatrix} x_0 + x_3 & x_1 + ix_2 \\ x_1 - ix_2 & x_0 - x_3 \end{pmatrix} \in \text{Herm}(2).$$

Under this identification one gets $\langle m, m \rangle = -\det(m)$ for all $m \in \text{Herm}(2)$. The complex Lie group $\text{SL}(2, \mathbb{C})$ acts naturally on $\mathbb{L}^4$ by $\Phi \cdot m = \Phi m \Phi^*$, being $\Phi \in \text{SL}(2, \mathbb{C}), \Phi^* = \Phi^t$, and $m \in \text{Herm}(2)$. Consequently, $\text{SL}(2, \mathbb{C})$ preserves the metric and the orientations. We shall view the hyperbolic 3-space of negative curvature $-r^2$ in its Minkowski model, that is, $\mathbb{H}^3(-r^2) = \{ x \in \mathbb{L}^4 : \langle x, x \rangle = -1/r^2, x_0 > 0 \}$. The above identification makes $\mathbb{H}^3(-r^2)$ become

$$\mathbb{H}^3(-r^2) = \left\{ \frac{1}{r} \Phi \Phi^* : \Phi \in \text{SL}(2, \mathbb{C}) \right\}, \quad (r > 0).$$

In the same way, the *de Sitter* space $\mathbb{S}^3_1(r^2) = \{ x \in \mathbb{L}^4 : \langle x, x \rangle = 1/r^2 \}$ is regarded as

$$\mathbb{S}^3_1(r^2) = \left\{ \frac{1}{r} \Phi \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \Phi^* : \Phi \in \text{SL}(2, \mathbb{C}) \right\}, \quad (r > 0).$$

We shall use the notation $\mathbb{H}^3 = \mathbb{H}^3(-1)$ and $\mathbb{S}^3_1 = \mathbb{S}^3_1(1)$.

Finally, the *positive light cone* $\mathbb{N}^3 = \{ x \in \mathbb{L}^4 : \langle x, x \rangle = 0, x_0 > 0 \}$ is seen as the space of positive semi-definite matrices in $\text{Herm}(2)$ with determinant 0, and can be described as

$$\mathbb{N}^3 = \{ w w^* : w^t = (w_1, w_2) \in \mathbb{C}^2 \setminus \{(0,0)\} \},$$

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where \( w \in \mathbb{C}^2 \setminus \{(0,0)\} \) is uniquely defined up to multiplication by an unimodular complex number. The quotient \( \mathbb{N}^3/\mathbb{R}^+ \) inherits a conformal structure and it can be regarded as the ideal boundary \( \mathbb{S}^2_{\infty} \) of the hyperbolic 3-space \( \mathbb{H}^3 \) in \( \mathbb{L}^4 \). The map \( \psi w \to [(w_1, w_2)] \) becomes the quotient map of \( \mathbb{N}^3 \) onto \( \mathbb{S}^2_{\infty} \) and identifies \( \mathbb{S}^2_{\infty} \) with \( \mathbb{C} \mathbb{P}^1 \equiv \mathbb{C} \cup \{ \infty \} \).

An immersion \( \psi : \Sigma \to \mathbb{L}^4 \) of a connected orientable surface \( \Sigma \) is said to be a spacelike surface if \( \Sigma \) inherits via \( \psi \) a Riemannian metric. Thus we shall regard \( \Sigma \) as a Riemann surface with the conformal structure induced by \( \psi \).

Let \( \psi : \Sigma \to \mathbb{L}^4 \) be a spacelike surface, and choose a local conformal coordinate \( z \) on \( \Sigma \) and an oriented orthonormal frame \( \{ \eta, \tilde{\eta} \} \) of \( T^\perp \Sigma \), being \( \tilde{\eta} \) a timelike vector field with values in \( \mathbb{H}^3 \), and \( \eta \) a spacelike one. Thus the induced metric of \( \Sigma \) is written as \( ds^2 = \lambda |dz|^2 \) for some positive smooth function \( \lambda \). If we define the moving frame

\[
\sigma = (\psi_z, \psi_{\bar{z}}, \eta, \tilde{\eta})^T
\]

the structure equations for the immersion are

\[
\sigma_z = U \sigma, \quad \sigma_{\bar{z}} = V \sigma,
\]

where

\[
U = \begin{pmatrix}
(\log \lambda)_z & 0 & p & \bar{p} \\
0 & 0 & E & \bar{E} \\
-2E/\lambda & -2p/\lambda & 0 & A \\
2\bar{E}/\lambda & 2\bar{p}/\lambda & A & 0 \\
\end{pmatrix}, \quad V = \begin{pmatrix}
0 & 0 & E & \bar{E} \\
0 & (\log \lambda)_{\bar{z}} & \bar{p} & \bar{p} \\
-2\bar{p}/\lambda & -2E/\lambda & 0 & \bar{A} \\
2\bar{p}/\lambda & 2\bar{E}/\lambda & \bar{A} & 0 \\
\end{pmatrix},
\]

and

\[
E = \langle \psi_z \bar{z}, \eta \rangle, \quad \bar{E} = -\langle \psi_{\bar{z}} z, \bar{\eta} \rangle, \quad A = -\langle \eta_z, \bar{\eta} \rangle, \\
p = \langle \psi_z z, \eta \rangle, \quad \bar{p} = -\langle \psi_{\bar{z}} \bar{z}, \bar{\eta} \rangle.
\]

The integrability condition for this system,

\[
U_z - V_{\bar{z}} + [U, V] = 0,
\]

turns into the following Gauss-Codazzi-Ricci equations:

\[
\begin{align*}
\text{Gauss:} & \quad (\log \lambda)_{zz} = \frac{2}{\chi} (|p|^2 - |\bar{p}|^2 + E^2 - \bar{E}^2), \\
\text{Codazzi (1):} & \quad p_z - E_z = A\bar{E} - \bar{A}\bar{p} - E(\log \lambda)_z, \\
& \quad \bar{p}_z - \bar{E}_z = AE + \bar{A}p - \bar{E}(\log \lambda)_z, \\
\text{Codazzi (2):} & \quad \left(\frac{\bar{q}}{\chi}\right)_z - \left(\frac{q}{\chi}\right)_{\bar{z}} = \frac{1}{\chi} (\bar{A}\bar{E} - \bar{A}\bar{p} - \bar{p}(\log \lambda)_z) \\
& \quad \left(\frac{\bar{q}}{\chi}\right)_{\bar{z}} - \left(\frac{q}{\chi}\right)_z = \frac{1}{\chi} (AE - Ap - p(\log \lambda)_z) \\
\text{Ricci:} & \quad A_{\bar{z}} - \bar{A}_z = -\frac{4i}{\chi} \text{Im}(p\bar{p}).
\end{align*}
\]
The mean curvature vector of the immersion \( \psi : \Sigma \to \mathbb{L}^4 \) will be denoted by \( H : \Sigma \to \mathbb{L}^4 \), and with the above notations it is given by
\[
H = \frac{2}{\lambda}(E\eta + \tilde{E}\tilde{\eta}).
\]

(2.5)

Besides, as \( \eta + \tilde{\eta} \) takes its values in the light cone \( \mathbb{N}^3 \), we may define on any spacelike surface in \( \mathbb{L}^4 \) the map \([\eta + \tilde{\eta}] : \Sigma \to S_{\infty}^2 \equiv \mathbb{C} \cup \{\infty\}\). It is straightforward to check that this map does not depend on the chosen orthonormal frame \( \{\eta, \tilde{\eta}\} \) of the oriented normal bundle. So, the following definition makes sense:

**Definition 1** The map \( G = [\eta + \tilde{\eta}] : \Sigma \to \mathbb{C} \cup \{\infty\} \) is called the hyperbolic Gauss map of the spacelike surface \( \psi : \Sigma \to \mathbb{L}^4 \).

The present paper deals with spacelike surfaces in \( \mathbb{L}^4 \) with isotropic mean curvature vector, that is, surfaces satisfying \( \langle H, H \rangle = 0 \). We shall call any such surface a *marginally trapped surface*. Observe that, with this definition, any spacelike surface with vanishing mean curvature in \( \mathbb{L}^4 \) is marginally trapped. After a change of orientation in the normal bundle if necessary (i.e. after a change of sign in \( \eta \)), the above condition is written as \( E = \tilde{E} \). Apart from their interest in Relativity Theory, the geometric importance of marginally trapped surfaces comes from the following fact.

**Lemma 2** Let \( \psi : \Sigma \to \mathbb{L}^4 \) be a marginally trapped surface. Then its hyperbolic Gauss map \( G : \Sigma \to \mathbb{C} \cup \{\infty\} \) is conformal.

This result follows simply by noting that
\[
\langle (\eta + \tilde{\eta})_z, (\eta + \tilde{\eta})_z \rangle = 4(E - \tilde{E})(p - \tilde{p}) = 0
\]
for every marginally trapped surface in \( \mathbb{L}^4 \).

**Remark 3** Let \( \psi : \Sigma \to \mathbb{L}^4 \) be a marginally trapped surface in \( \mathbb{L}^4 \) that actually lies in some \( \mathbb{R}^3 \subset \mathbb{L}^4 \), \( \mathbb{L}^3 \subset \mathbb{L}^4 \), \( \mathbb{H}^3(-r^2) \subset \mathbb{L}^4 \) or \( S_1^3(r^2) \subset \mathbb{L}^4 \). Then, by a straightforward computation, the condition \( \langle H, H \rangle = 0 \) implies that \( \psi \) has zero mean curvature if it lies in some \( \mathbb{R}^3 \) or some \( \mathbb{L}^3 \). In the same way, if \( \psi \) lies in some \( \mathbb{H}^3(-r^2) \) or some \( S_1^3(r^2) \), its mean curvature in that ambient space is constant, of value \( H = r \). In all these cases, the mean curvature vector \( H \) of \( \psi \) in \( \mathbb{L}^4 \) is parallel.

From now on we shall work with marginally trapped surfaces with *flat normal bundle*, that is, we shall assume that the normal curvature vanishes identically, \( R^4 \equiv 0 \). This amounts to say that \( A_z - \overline{A}_z = 0 \). But it comes clear that this condition implies the local existence of a real function \( \beta \) on \( \Sigma \) such that \( d\beta = Adz + \tilde{A}\tilde{d}z \). Thus, by considering the new normal frame given by
\[
\xi = \cosh(\beta)\eta - \sinh(\beta)\tilde{\eta}, \quad \tilde{\xi} = -\sinh(\beta)\eta + \cosh(\beta)\tilde{\eta}
\]
we can assume that \( A = 0 \) holds in the structure equations (2.2). In other words, there exists an orthonormal frame \( \xi, \tilde{\xi} \) of the normal bundle that is parallel. Let us also remark
that this parallel orthonormal frame is unique up to constant hyperbolic rotations in
the Lorentzian normal bundle of the immersion.

We shall keep denoting by \( \{ \eta, \tilde{\eta} \} \) the new parallel orthonormal frame \( \{ \xi, \tilde{\xi} \} \) of \( T^1\Sigma \). It is not difficult to see that with the above hypothesis the equations (2.2) and (2.4) can be simplified considerably.

**Lemma 4** Let \( \psi : \Sigma \to \mathbb{L}^4 \) be a marginally trapped surface with flat normal bundle, and let \( z \) denote a local conformal coordinate on \( \Sigma \). Then there is an orthonormal frame \( \{ \eta, \tilde{\eta} \} \) of \( T^1\Sigma \) such that the moving frame (2.1) satisfies (2.2) for the matrices

\[
U = \begin{pmatrix}
\log \lambda & 0 & \bar{p} \\
0 & 0 & E & E \\
-2E/\lambda & -2\bar{p}/\lambda & 0 & 0 \\
2E/\lambda & 2\bar{p}/\lambda & 0 & 0
\end{pmatrix}, \quad V = \begin{pmatrix}
0 & 0 & E & E \\
0 & (\log \lambda) \bar{z} & \bar{p} & \tilde{p} \\
-2\bar{p}/\lambda & -2E/\lambda & 0 & 0 \\
2\bar{p}/\lambda & 2E/\lambda & 0 & 0
\end{pmatrix}.
\]

Here \( E, p, \tilde{p} \) are as in (2.3), and they obey the integrability conditions

\[
\text{Gauss-Ricci:} \quad (\log \lambda)_{z\bar{z}} = \frac{2(|p|^2 - |\bar{p}|^2)}{\lambda} = \frac{2\bar{p}(p - \bar{p})(p + \bar{p})}{\lambda}, \quad (2.7)
\]

\[
\text{Codazzi:} \quad p_{\bar{z}} = \bar{p}_{\bar{z}} = \lambda \left( \frac{E}{\lambda} \right)_{\bar{z}}.
\]

The mean curvature vector of any surface in the conditions of the above Lemma is given by

\[
H = \frac{2E}{\lambda} (\eta + \tilde{\eta}) \quad (2.8)
\]

It is immediate from the Codazzi equations that \( H \) is parallel if and only if \( E/\lambda \) is constant, if and only if both \( p, \bar{p} \) are holomorphic.

Besides, let us note this other consequence of the Codazzi equations.

**Lemma 5** The quadratic differential \( Q = (\bar{p} - p)dz^2 \) is holomorphic on every marginally trapped surface with flat normal bundle in \( \mathbb{L}^4 \).

From now on we will denote the holomorphic function \( \bar{p} - p \) as \( q = \bar{p} - p \). We shall call \( Q = q(z)dz^2 \) the Hopf differential of the surface.

**Remark 6** If \( \psi : \Sigma \to \mathbb{L}^4 \) is a marginally trapped surface with flat normal bundle on which \( Q \) vanishes identically, then by (2.6), \( \eta + \tilde{\eta} \) is constant and \( \Sigma \) is flat. Thus \( \psi(\Sigma) \) lies in an affine degenerate hyperplane of \( \mathbb{L}^4 \). Flat surfaces lying in degenerate hyperplanes of \( \mathbb{L}^4 \) were completely described in explicit coordinates in [GMM2]. Let us also indicate that many of these surfaces are complete.

In the remaining of the present work we will assume that \( Q \) does not vanish identically.
3 A conformal representation

Let us start this section recalling a classical result by Ricci, stating that a necessary and sufficient condition for a Riemannian surface $(S, ds^2)$ to be locally isometric to a minimal surface in $\mathbb{R}^3$ is that it has non-positive curvature $K \leq 0$ and the conformal pseudo-metric $\tilde{ds}^2 = \sqrt{-K} ds^2$ is flat in case $K \neq 0$. This is equivalent to impose that the conformal pseudo-metric $-K ds^2$ has constant curvature 1 at its regular points.

Analogously, it can be easily proved that $(S, ds^2)$ is locally isometric to a maximal surface in $\mathbb{R}^3$ if and only if $K \geq 0$ and $\sqrt{K} ds^2$ is flat, if and only if $K \geq 0$ and $K ds^2$ has constant curvature $-1$.

**Definition 7** Let $\psi : \Sigma \to \mathbb{L}^4$ be a marginally trapped surface with flat normal bundle. Then $\psi$ is said to be a surface of Bryant type in $\mathbb{L}^4$ provided $\Sigma$ is locally isometric to some minimal surface in $\mathbb{R}^3$ or to some maximal surface in $\mathbb{L}^3$.

Let us examine this kind of surfaces. For this we start with a simply connected marginally trapped surface with flat normal bundle $\psi : \Sigma \to \mathbb{L}^4$. From the Gauss equation in (2.7) we find that its Gauss curvature is

$$K = \frac{4 (|\tilde{p}|^2 - |p|^2)}{\lambda^2}.$$

We shall denote $\varepsilon = \text{sign}(K) = \text{sign} (|\tilde{p}|^2 - |p|^2)$. If $\varepsilon = 0$ the surface is of Bryant type. Otherwise we can define on $\Sigma$ away from the isolated flat points of $\Sigma$ the metric $\tilde{d}s^2 = \sqrt{\varepsilon K} ds^2$, and the above formula easily implies that $\tilde{d}s^2$ is flat if and only if

$$(\log (\varepsilon(|\tilde{p}|^2 - |p|^2)))_{\tilde{z}z} = 0,$$

that is, if and only if there is a holomorphic function $\varphi : \Sigma \to \mathbb{C}$ such that $\varepsilon(|\tilde{p}|^2 - |p|^2) = |\varphi|^2$ (recall that $\Sigma$ is simply connected). But by the Ricci equation, this is equivalent to the fact that

$$\tilde{p} + p = \varepsilon |f|^2 q,$$ (3.1)

being

$$f = \varphi/q : \Sigma \to \mathbb{C} \cup \{\infty\}$$ (3.2)

a meromorphic function on $\Sigma$.

**Remark 8** As we know that a marginally trapped surface with flat normal bundle has parallel mean curvature if and only if both $p, \tilde{p}$ are holomorphic, it turns out that a surface of Bryant type has parallel mean curvature if and only if $f$ is constant.

Moreover, if $\psi : \Sigma \to \mathbb{L}^4$ is a marginally trapped surface with flat normal bundle and parallel mean curvature vector in $\mathbb{L}^4$, then it belongs to a known family. Indeed, if this is the case we get from the Ricci equation that $\tilde{p} = kp$, $k \in \mathbb{R}$, which, jointly with the structure equations (2.6) provide $\psi_z = k'(k\eta + \tilde{\eta})_z$, $k' \in \mathbb{R}$. This identity implies that $\psi$ must lie in a totally umbilical affine hyperquadric of $\mathbb{L}^4$, and thus is a known example by Remark 3.
The basic result of the present work is a conformal representation for surfaces of Bryant type in $\mathbb{L}^4$.

**Theorem 9 (Conformal representation)** Let $\psi : \Sigma \to \mathbb{L}^4$ be a non-flat simply connected surface of Bryant type, and let $\varepsilon = \pm 1$ denote the sign of its Gaussian curvature. Then there exist a meromorphic function $g$ and a holomorphic 1-form $\omega$ on $\Sigma$, and three constants $a, b \in \mathbb{R}$, $c \in \mathbb{C}$ satisfying

\begin{equation}
C.1 \quad 1 - \varepsilon |g|^2 > 0, \text{ and the zeros of } \omega \text{ of order } 2k \text{ correspond to the poles of } g \text{ of order } k, \text{ and}
\end{equation}

\begin{equation}
C.2 \quad \text{There is a holomorphic function } f \text{ verifying that}
\end{equation}

$$df = \left(c + (a + \varepsilon b)g + \varepsilon \bar{c}g^2\right) \omega$$

and that $\omega df/f$ is holomorphic,

such that the immersion can be expressed as

$$\psi = F\Omega F^* : \Sigma \to \mathbb{L}^4.$$  

(3.4)

Here $F : \Sigma \to \text{SL}(2, \mathbb{C})$ is a meromorphic curve satisfying

$$F^{-1}dF = \begin{pmatrix} 0 & (a + \varepsilon \bar{c})\omega \\ dg/f & 0 \end{pmatrix} =: \mathcal{A}$$

(3.5)

and $\Omega : \Sigma \to \text{Herm}(2)$ verifies the differential equation

$$d\Omega + \mathcal{A}\Omega + \Omega^* = \begin{pmatrix} 2\varepsilon \text{Re} (g\bar{f}\omega) & (1 - \varepsilon |g|^2) \omega \\ (1 - \varepsilon |g|^2) \overline{\omega} & 0 \end{pmatrix}.$$  

(3.6)

Conversely, let $\Sigma$ be a simply connected Riemann surface, $\varepsilon = \pm 1$, and consider a meromorphic function $g$ and a holomorphic 1-form $\omega$ on $\Sigma$ satisfying $C.1$ and $C.2$ for some constants $a, b \in \mathbb{R}$, $c \in \mathbb{C}$. Then there exist a meromorphic curve $F : \Sigma \to \text{SL}(2, \mathbb{C})$ satisfying (3.5), and a solution $\Omega : \Sigma \to \text{Herm}(2)$ to the system (3.6). Moreover, the map $\psi : \Sigma \to \mathbb{L}^4$ given by (3.4) is a surface of Bryant type in $\mathbb{L}^4$ for which $\varepsilon$ is the sign of its Gauss curvature.

**Remark 10** The only surfaces of Bryant type in $\mathbb{L}^4$ that are flat are those in Remark 6. So, it is not restrictive to assume in the representation theorem that the surfaces of Bryant type are non-flat.

**Proof**: Let $\psi : \Sigma \to \mathbb{L}^4$ be a non-flat simply connected surface of Bryant type. As $Q$ is a non-zero holomorphic 2-form, $\Sigma$ cannot be the Riemann sphere, and so we may choose a global holomorphic coordinate $z$ on the Riemann surface $\Sigma$. Following the notations of Lemma 4 we get that $q = \bar{p} - p$ is holomorphic and $f$ as in (3.2) is meromorphic. On the other hand it is easy to check that $d\sigma^2 = \varepsilon Kds^2$ is a pseudometric on $\Sigma$ of
constant curvature $-\varepsilon$. As $\Sigma$ is simply connected, by the Frobenius theorem there exists a meromorphic function $g$ (holomorphic with $|g| < 1$ if $\varepsilon = 1$) on $\Sigma$ such that (see [Bry, GMM3, GaMi])

$$\varepsilon K ds^2 = \frac{4|dg|^2}{(1 - \varepsilon|g|^2)^2}.$$ 

Now, since from (3.1) and (3.2) we know that

$$K = \varepsilon \frac{4|f|^2|q|^2}{\lambda^2},$$

it follows that

$$\frac{|f|^2|q|^2}{\lambda} = \frac{|g_z|^2}{(1 - \varepsilon|g|^2)^2}. \quad (3.7)$$

Thus

$$\lambda = \left| \frac{f q}{g_z} \right|^2 \left(1 - \varepsilon|g|^2 \right)^2. \quad (3.8)$$

In this way $\omega = fQ/dg$ is a meromorphic 1-form on $\Sigma$, and it has no poles. Note that $Q = \omega dg/f$, so this quantity defines a holomorphic quadratic differential on $\Sigma$. Besides, the zeros of $\omega$ of order $2k$ must trivially coincide with the poles of $g$ of order $k$.

Since from Lemma 2 we know that the hyperbolic Gauss map $G = [\eta + \bar{\eta}] : \Sigma \to S^2_{\infty}$ is conformal, there exist holomorphic functions $A, B : \Sigma \to \mathbb{C}$ and a positive real function $\mu : \Sigma \to \mathbb{R}^+$ such that

$$\eta + \bar{\eta} = \mu \begin{pmatrix} A\bar{A} & A\bar{B} \\ \bar{A}B & \bar{B}B \end{pmatrix}. \quad (3.9)$$

Thus $\langle (\eta + \bar{\eta})_z, (\eta + \bar{\eta})_{\bar{z}} \rangle = \frac{1}{2} |AB_z - BA_{\bar{z}}|^2 \mu^2$, and since from (2.1) and (2.6) it is obtained

$$(\eta + \bar{\eta})_{\bar{z}} = \frac{2\nu}{\lambda} \psi_z, \quad (3.10)$$

we get by means of (3.8)

$$|ADB - BdA|^2 \mu^2 = \frac{4|dg|^2}{|f|^2 (1 - \varepsilon|g|^2)^2}. \quad (3.11)$$

Besides, from (3.11) we see that

$$\left| \frac{dg}{f(ADB - BdA)} \right| = \frac{1}{2} \mu(1 - \varepsilon|g|^2) > 0.$$ 

So, $dg/(f(ADB - BdA))$ never vanishes, and all its poles are of even order. This ensures the existence of a meromorphic function (holomorphic if $\varepsilon \neq -1$) $S$ verifying

$$S^2 = \frac{dg}{f(ADB - BdA)}.$$
If we now choose $C = AS$, $D = BS$ we find that $[(A, B)] = [(C, D)]$ and $CdD - DdC = dg/f$. Thus, by substituting $(A, B)$ with $(C, D)$ and $\mu$ with $\varrho$ so that

$$\eta + \tilde{\eta} = \varrho \left( \begin{array}{cc} CC & CD \\ CD & DD \end{array} \right),$$

equation (3.11) turns into

$$\varrho = \frac{2}{1 - \varepsilon |g|^2}. \tag{3.12}$$

Now consider the meromorphic curve $F : \Sigma \to \text{SL}(2, \mathbb{C})$

$$F = \left( \begin{array}{cc} C & \frac{f dC/dg}{dG/f} \\ D & \frac{f dD/dg}{dG/f} \end{array} \right). \tag{3.13}$$

Then there exists a meromorphic 1-form $\vartheta$ on $\Sigma$ such that

$$F^{-1}dF = \left( \begin{array}{cc} 0 & \vartheta \\ dG/f & 0 \end{array} \right). \tag{3.14}$$

Moreover,

$$\eta + \tilde{\eta} = F \left( \begin{array}{cc} \varrho & 0 \\ 0 & 0 \end{array} \right) F^* \tag{3.15}$$

and from there, (3.14) and (3.12), we have

$$(\eta + \tilde{\eta})_z = F \begin{pmatrix} \frac{2\varepsilon gg_z}{(1 - \varepsilon |g|^2)^2} & \frac{2g_z}{f(1 - \varepsilon |g|^2)} \\ 0 & 0 \end{pmatrix} F^*. \tag{3.16}$$

Once here we recall (3.10) and the fact that $q$ does not vanish identically to obtain from the above expression that

$$\psi_z = \frac{\lambda}{2\eta} F \begin{pmatrix} \frac{2\varepsilon gg_z}{(1 - \varepsilon |g|^2)^2} & \frac{2g_z}{f(1 - \varepsilon |g|^2)} \\ 0 & 0 \end{pmatrix} F^*. \tag{3.16}$$

In addition, using (3.7) we obtain the final expression for $\psi_z$,

$$\psi_z = F \begin{pmatrix} \varepsilon |f|^2 g g_z & (1 - \varepsilon |g|^2) q f \\ g_z & g_z \end{pmatrix} F^*. \tag{3.16}$$

Finally, let us note that as $\text{SL}(2, \mathbb{C})$ acts on $\mathbb{L}^4$ as the connected component of the identity in its isometry group, the immersion $\psi : \Sigma \to \mathbb{L}^4$ can be expressed as $\psi = F\Omega F^* : \Sigma \to \mathbb{L}^4$ for an adequately chosen intermediate matrix $\Omega : \Sigma \to \text{Herm}(2)$. It comes clear from (3.16) that $\Omega$ is a solution of the differential system (3.6). Next, note
that by differentiation of (3.16) with respect to $\bar{z}$ and noting that $\psi_{\bar{z}}$ is real, we see that the data $f, g, \omega, \vartheta$ must verify

$$
\varepsilon g \omega \left( df - \bar{g} \vartheta \right) + \omega \vartheta = \varepsilon \bar{g} \bar{\omega} \left( df - g \vartheta \right) + \bar{\omega} \vartheta.
$$

Thus, by Lemma 16 in the Appendix, and as $g$ is not constant (otherwise $\psi$ would be flat), we obtain the existence of constants $a, b \in \mathbb{R}$ and $c \in \mathbb{C}$ such that

$$
\vartheta = (a + \varepsilon \bar{c}g)\omega, \quad df = (c + (a + \varepsilon b)g + \varepsilon \bar{c}g^2) \omega.
$$

(3.17)

Particularly, $f, \vartheta$ are holomorphic, and C.2 and (3.5) hold. Thus the proof of the first part of the theorem is complete.

For the converse, we start with the Weierstrass data $(g, \omega)$ and constants $a, b \in \mathbb{R}, c \in \mathbb{C}$ verifying C.1 and C.2 on a simply connected Riemann surface $\Sigma$. Then the system (3.5) has a (possibly multivalued) solution $F : \Sigma \to \text{SL}(2, \mathbb{C})$ of the form (3.13), where $C, D$ are linearly independent solutions of (see [GMM1])

$$
Z'' - \frac{(g'/f)'}{g'/f}Z' - (a + \varepsilon \bar{c}g)qZ = 0, \quad \left( ' = \frac{d}{dz} \right),
$$

(3.18)

being $Q = q(z)dz^2 = \omega dg/f$.

As the meromorphic 1-form in (3.5) has its poles at the poles of $g$, we see that the solution $F$ is locally well defined and holomorphic away from the poles of $g$.

Let now $z_0 \in \Sigma$ be a pole of $g$ of order $k \geq 1$, and let $\delta \geq 0$ denote the order of the zero of $f$ at $z_0$ (possibly $\delta = 0$). From C.2 it is clearly seen that if $c \neq 0$, then $\delta \in \{0, 1\}$, while if $c = 0$ then $\delta = 0$ (otherwise we would have $\delta = k + 1$ by C.2, which would contradict that $\omega dg/f$ has no poles).

It is straightforward that $(g'/f)'/(g'/f)$ has a simple pole at $z_0$, of residue $-(k+\delta+1)$. Besides, if $h = (a + \varepsilon \bar{c}g)q$, it follows directly that $h$ is holomorphic at $z_0$ if $c = 0$, it has a simple pole at $z_0$ if $c \neq 0$ and $f(z_0) \neq 0$, and has a pole of order two at $z_0$ otherwise. From this, a simple calculation shows that

$$
h_{-2} = \lim_{z \to z_0} (z - z_0)^2 h(z) = \delta k.
$$

Hence, the differential equation (3.18) has a regular singularity at $z_0$ (see [CoLe]). Its indicial equation is

$$
\lambda^2 + (k + \delta)\lambda + k\delta = 0,
$$

that has the integer roots $-k$ and $-\delta$. Therefore, both $C, D$ are single valued meromorphic functions on $\Sigma$, and the orders of their poles at $z_0$ are $k$ and $\delta$ (see [GMM1]). Particularly, any pole of $C$ or $D$ of order $l$ is located at a poles of $g$ of order $\geq l$.

Once here, we have ensured the existence of a meromorphic solution $F : \Sigma \to \text{SL}(2, \mathbb{C})$ to (3.5) of the form (3.13). If we now set $\varrho : \Sigma \to [0, +\infty)$ as (3.12), the map

$$
N = F\left( \begin{array}{cc} \varrho & 0 \\ 0 & 0 \end{array} \right) F^*
$$

(3.19)
has a finite value at every point, due to the previous analysis regarding the poles of $C, D$. Thus we have a map $N : \Sigma \to \mathbb{N}^3$. The same argument shows that the 1-forms

$$\phi dz = F \left( \begin{array}{cc} \varepsilon g f \omega & (1 - \varepsilon|g|^2) \omega \\ 0 & 0 \end{array} \right) F^*, \quad \tilde{\phi}d\tilde{z} = F \left( \begin{array}{cc} \varepsilon \tilde{g} f \tilde{\omega} & 0 \\ (1 - \varepsilon|g|^2) \tilde{\omega} & 0 \end{array} \right) F^*$$

take finite values at all points. But now, noting that $\phi^* = \tilde{\phi}$ and that $\phi_z = \tilde{\phi}_z$, we can conclude the existence of a map $\psi : \Sigma \to \text{Herm}(2)$ such that $\psi_z = \phi$ and $\psi_{\bar{z}} = \phi$. Finally, let us define $\Omega : \Sigma \to \text{Herm}(2)$ as $\Omega = F^{-1} \psi(F^{-1})^*$, whose entries may take infinite values at some points. Then $\psi = F\Omega F^*$ and $\Omega$ is trivially a solution to the differential system (3.6).

At last, as by differentiation of (3.16) we have that $\psi_{z\bar{z}}$ is collinear with $N$, we conclude that $\psi$ is a marginally trapped surface. Now, deriving the right hand side of (3.19) with respect to $z$ we get that $N$ is parallel, that is $\psi$ has flat normal bundle. And as $\langle d\psi, d\psi \rangle = (1 - \varepsilon|g|^2)^2|\omega|^2$, $\psi$ is regular and of Bryant type, and the proof is finished.

\[ \bullet \]

**Remark 11** The mean curvature vector of the Bryant-type surface $\psi : \Sigma \to \mathbb{L}^4$ verifies the relation $2\psi_{z\bar{z}} = \lambda H = 2E(\eta + \bar{\eta})$. This indicates by differentiation of (3.16) that

\[ \frac{2E}{\lambda} = \frac{a + b|g|^2 + 2\varepsilon \text{Re}(\bar{c}g)}{1 - \varepsilon|g|^2}. \quad (3.20) \]

Thus an explicit expression for $H$ in terms of the Weierstrass data is obtained from (2.8), (3.20) and (3.15).

**Remark 12** The hyperbolic Gauss map $G : \Sigma \to \mathbb{C} \cup \{\infty\}$ of a surface of Bryant type is a geometric concept, and thus is uniquely determined at every point. However, this is not the case for the other basic meromorphic data of a Bryant surface. First, note that as $Q$ depends on the chosen frame of the normal bundle, it is defined up to the change $Q \to e^{\alpha}Q$, where $\alpha \in \mathbb{R}$ is the constant hyperbolic angle relating two different frames. Therefore, the function $f$ is by definition defined up to $f \to e^{-\alpha + i\beta} f$, $\beta \in \mathbb{R}$. Besides, the meromorphic function $g$ is unique up to isometries of the 2-sphere $S^2 \equiv \mathbb{C} \cup \{\infty\}$ if $\varepsilon = -1$, and up to isometries of the Poincaré disk $\mathbb{H}^2 \equiv \mathbb{D}$ if $\varepsilon = 1$. Thus, $g$ is unique up the change

$$g \to \frac{\tau g + \varepsilon \gamma}{\gamma g + \tau}, \quad |\tau|^2 - |\varepsilon| \gamma|^2 = 1.$$

Noting now that $Q = \omega dg/f$, the above comments show that $\omega$ is defined up to

$$\omega \to e^{i\beta}(\gamma g + \tau)^2 \omega.$$

To close this section, we shall relate the hyperbolic Gauss map $G : \Sigma \to \mathbb{C} \cup \{\infty\}$ to the Weierstrass data of a Bryant-type surface in $\mathbb{L}^4$. First, observe that in (3.9) we may
choose $A = 1$, $B = G$. Then, with these choices we end up with the formula $\psi = F\Omega F^*$ where
\[
F = \begin{pmatrix} C & fdC/dg \\ D & fdD/dg \end{pmatrix}, \quad C = \sqrt{dg/(fdG)}, \quad D = G\sqrt{g/(fdG)}.
\]
This formula extends a result by Small [Sma] for mean curvature one surfaces in $\mathbb{H}^3$ (see also [GMM3]).

An alternative relation between $G$ and the Weierstrass data relies in the concept of Schwarzian derivative \{h, z\} of a meromorphic function $h$:
\[
\{h, z\} = \left( \frac{h''}{h'} \right)' - \frac{1}{2} \left( \frac{h''}{h'} \right)^2, \quad \left( ' = \frac{d}{dz} \right).
\]
We get then by [GMM1, Eq. (28)] and (3.14) that the following relation holds on any surface of Bryant type in $\mathbb{L}^4$:
\[
\{G, z\}dz^2 = \left( \left( \frac{(g'/f)'}{g'/f} \right)' - \frac{1}{2} \left( \frac{(g'/f)'}{g'/f} \right)^2 \right) dz^2 - (a + \varepsilon\bar{c}g)Q.
\]
We remark that this formula extends an important equation due to Umehara and Yamada [UmYa2] in the context of mean curvature one surfaces in $\mathbb{H}^3$.

4 Examples

**Representation of CMC-$r$ surfaces:** next, we show that the conformal representation in Theorem 9 generalizes the Bryant representation [Bry] for surfaces with $H = r$ in $\mathbb{H}^3(-r^2)$, as well as the Aiyama-Akutagawa one [AiAk] for spacelike surfaces with $H = r$ in $S^3_1(r^2)$. For this, we shall use the unified notation $\mathbb{M}^3(\varepsilon r^2)$ to denote $\mathbb{H}^3(-r^2)$ for $\varepsilon = -1$ and $S^3_1(r^2)$ for $\varepsilon = 1$.

Let $\psi : \Sigma \to \mathbb{M}^3(\varepsilon r^2) \subset \mathbb{L}^4$ be a simply connected (spacelike) CMC-$r$ surface, and let $\eta_1$ be its unit normal in $\mathbb{M}^3(\varepsilon r^2)$. Then \{$\eta_1, \eta_2 := r\psi$\} is a parallel orthonormal frame in the normal bundle of $\psi$ in $\mathbb{L}^4$. So, using the notations of the first two sections, it follows directly that $\varepsilon Q = \langle \psi_{zz}, \eta_1 \rangle dz^2$. Therefore $|f| = 1$ on $\Sigma$ and as $f$ is defined up to constant rotations (note that in this case we are working with a uniquely determined frame in the normal bundle), we may assume that $f = 1$.

With all of this, the differential system (3.6) can be explicitly solved under the condition $\det(\Omega) = -\varepsilon/r^2$, to obtain
\[
\Omega = \frac{1}{r} \begin{pmatrix} -\varepsilon & \varepsilon \bar{g} \\ \varepsilon g & 1 - \varepsilon |g|^2 \end{pmatrix} = \frac{1}{r} \begin{pmatrix} 0 & i \\ i & -ig \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -\varepsilon \end{pmatrix} \begin{pmatrix} 0 & -i \\ -i & ig \end{pmatrix}.
\]
Finally, we derive with respect to $z$ the identity $\psi = F\Omega F^*$ and compare it with (3.16) to deduce that $\vartheta = r\omega$ for the holomorphic 1-form $\vartheta$ in (3.17). In conclusion, the Bryant surface $\psi : \Sigma \to \mathbb{M}^3(\varepsilon r^2)$ is recovered as
\[
\psi = \frac{1}{r} \mathcal{B} \left( \begin{pmatrix} 1 & 0 \\ 0 & -\varepsilon \end{pmatrix} \right) \mathcal{B}^*, \quad \mathcal{B} = F \left( \begin{pmatrix} 0 & i \\ i & -ig \end{pmatrix} \right).
\]
Here $\mathcal{B}: \Sigma \to \text{SL}(2, \mathbb{C})$ is a null holomorphic curve (i.e. $\det(d\mathcal{B}) = 0$), and $F: \Sigma \to \text{SL}(2, \mathbb{C})$ verifies (3.5) for $a = r$, $c = 0$. Thus, the Bryant representation in [Bry] and the Bryant-type representation in [AiAk] are recovered.

The Weierstrass representation: now, we prove that the classical Weierstrass representation for minimal surfaces in $\mathbb{R}^3$ and its analogue for maximal surfaces in $\mathbb{L}^3$ are also included in Theorem 9. In order to do so, we fix the notation $\mathbb{R}^3_\varepsilon$ to denote $\mathbb{R}^3 \equiv x_0 = 0 \subset \mathbb{L}^4$ if $\varepsilon = -1$, and $\mathbb{L}^3 \equiv x_3 = 0 \subset \mathbb{L}^4$ if $\varepsilon = 1$.

Let $\psi: \Sigma \to \mathbb{R}^3_\varepsilon \subset \mathbb{L}^4$ be a minimal (or maximal) surface in $\mathbb{R}^3_\varepsilon$, with unit normal $\eta_1$. Then $\{\eta_1, \eta_2 = \frac{1}{2}(1-\varepsilon, 0, 0, 1+\varepsilon)\}$ is a parallel orthonormal frame of the normal bundle of $\psi$ in $\mathbb{L}^4$, and arguing as above we get that $f = 1$. Moreover, as $H = 0$, by (3.20) we have $a = b = c = 0$. Thus the differential equation (3.5) can be explicitly integrated, and we obtain a solution as

$$F = \begin{pmatrix} 1 & 0 \\ g & 1 \end{pmatrix}: \Sigma \to \text{SL}(2, \mathbb{C}).$$

(4.1)

Let $\Omega: \Sigma \to \text{Herm}(2)$ be the solution to (3.6), which is in this case as

$$\Omega_z + \begin{pmatrix} 0 & 0 \\ g_z & 0 \end{pmatrix} \Omega = \begin{pmatrix} \varepsilon g \omega & (1-\varepsilon|g|^2)\omega \\ 0 & 0 \end{pmatrix}.$$  

(4.2)

If we write

$$\Omega = \begin{pmatrix} U & V \\ \bar{V} & W \end{pmatrix},$$

then we see from (4.2) that

$$U = 2\varepsilon \text{Re} \int g \omega, \quad (Ug + \bar{V})_z = \varepsilon g^2 \omega, \quad (U\bar{g} + V)_z = \omega.$$

Now, as $\psi(\Sigma) \subset \mathbb{R}^3_\varepsilon$ it holds

$$\psi = F\Omega F^* = \begin{pmatrix} U & U\bar{g} + V \\ Ug + \bar{V} & \varepsilon U \end{pmatrix}.$$

So, by the above computations we finally obtain the Weierstrass representation:

$$\psi = \text{Re} \int \left( (1+\varepsilon)g, 1+\varepsilon g^2, -i(1-\varepsilon g^2), (1-\varepsilon)g \right) \omega.$$

Analytic deformation of surfaces: we shall show now that, with the above notations, CMC-$r$ surfaces in $\mathbb{M}^3(\varepsilon r^2)$ can be analytically and isometrically deformed to minimal (or maximal) surfaces in $\mathbb{R}^3_\varepsilon$ as $r$ tends to zero. This result was obtained in the case $\varepsilon = -1$ by Umehara and Yamada [UmYa].

To do so, we consider first the translated spaces

$$\tilde{\mathbb{M}}^3(\varepsilon r^2) = \left\{ p - \frac{1}{2r}(1-\varepsilon, 0, 0, -1-\varepsilon) : p \in \mathbb{M}^3(\varepsilon r^2) \right\},$$

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where we will assume that \( M^3(r^2) = \{ x \in S^3_1 : x_3 < 0 \} \) is a de Sitter half-space. Our deformation process will rely on the fact that, as \( r \to 0 \), the spaces \( \tilde{M}^3(\varepsilon r^2) \) converge to \( \mathbb{R}^3_\varepsilon \).

Let \( \psi_{r_0} : \Sigma \to \tilde{M}^3(\varepsilon r_0^2) \) be a simply connected CMC-\( r_0 \) surface, and choose \( z_0 \in \Sigma \). Let now \( (g, \omega) \) be its Weierstrass data, and suppose without loss of generality that \( g(z_0) = 0 \). Then, for every \( r > 0 \) there exists a unique (up to rigid motions) CMC-\( r \) surface in \( \tilde{M}^3(\varepsilon r^2) \) \( \psi_r : \Sigma \to \tilde{M}^3(\varepsilon r^2) \) that has the same Weierstrass data \( (g, \omega) \), and so it is isometric to the original immersion \( \psi_{r_0} \). Now, we know that \( \psi_r = \frac{1}{r} F_r \Delta F^*_r \), where \( F_r \in \text{SL}(2, \mathbb{C}) \),

\[
F_r^{-1} dF_r = \begin{pmatrix} 0 & r \omega \\ d\bar{g} & 0 \end{pmatrix}, \quad F_r(z_0) = \text{Id}, \quad \Delta = \begin{pmatrix} -\varepsilon & \varepsilon \bar{g} \\ \varepsilon g & 1 - \varepsilon |g|^2 \end{pmatrix}.
\]

Consider now the translated immersions \( X_r : \Sigma \to \tilde{M}^3(\varepsilon r^2) \) given by

\[
X_r = \psi_r - \frac{1}{2r} (1 - \varepsilon, 0, 0, -1 - \varepsilon) = \frac{1}{r} \left( F_r \Delta F^*_r - \begin{pmatrix} -\varepsilon & 0 \\ 0 & 1 \end{pmatrix} \right).
\]

As the family \( A_r := F_r^{-1} dF_r \) is real analytic with respect to \( r \in \mathbb{R} \), the family \( F_r : \Sigma \to \text{SL}(2, \mathbb{C}) \) is also real analytic with respect to \( r \). So, using that \( F_0 \) is given by (4.1), it is easy to see that

\[
F_r \Delta F^*_r - \begin{pmatrix} -\varepsilon & 0 \\ 0 & 1 \end{pmatrix} = a_1(z, \bar{z}) r + o(r).
\]

This assures that the family of surfaces \( X_r : \Sigma \to \mathbb{L}^4 \) is real analytic with respect to \( r \in \mathbb{R} \). Furthermore, \( X_0 : \Sigma \to \mathbb{R}^3_\varepsilon \) is given by

\[
X_0 = a_1(z, \bar{z}) = \frac{\partial}{\partial r} \bigg|_{r=0} F_r \Delta F^*_r.
\]

Finally, \( X_0 \) has zero mean curvature in \( \mathbb{R}^3_\varepsilon \). This happens because, by

\[
\frac{\partial^2}{\partial z \partial \bar{z}} (F_r \Delta F^*_r) = F_r \left( r^2 |\omega|^2 (1 - \varepsilon |g|^2) \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \right) F^*_r
\]

and the analyticity of the family in \( r \in \mathbb{R} \), we have

\[
\frac{\partial^2 X_0}{\partial z \partial \bar{z}} = \frac{\partial^2}{\partial z \partial \bar{z}} \left( \frac{\partial}{\partial r} \bigg|_{r=0} F_r \Delta F^*_r \right) = \frac{\partial}{\partial r} \bigg|_{r=0} \left( \frac{\partial^2}{\partial z \partial \bar{z}} (F_r \Delta F^*_r) \right) = 0.
\]

Therefore, we conclude that the CMC-\( r \) surfaces of \( \mathbb{H}^3(-r^2) \) (resp. \( S^3_1(r^2) \)) can be perturbed in an analytic and isometric way to minimal surfaces in \( \mathbb{R}^3 \) (resp. to maximal surfaces in \( \mathbb{L}^3 \)) as \( r \) approaches to zero.

**New families of complete examples:** there exist many complete surfaces of Bryant type in \( \mathbb{L}^4 \) with non-parallel mean curvature, that can be constructed by means of the representation formula.
To see this, let us choose $c = 0$ in Theorem 9. Then $df = (a + \varepsilon b)g\omega$. Therefore, the condition C.2 in this case just asks for the existence of a nowhere-zero primitive $\int g\omega$ of the holomorphic 1-form $g\omega$ on $\Sigma$, a condition that always holds locally. Moreover, if $a + \varepsilon b \neq 0$ the example has non-parallel mean curvature, and if $a = -\varepsilon b \neq 0$ the surface lies in some affine hyperbolic 3-space or de Sitter 3-space in $\mathbb{L}^4$ and is a Bryant surface there (see Remark 8). Finally, if $a = -\varepsilon b = 0$, the surface has zero mean curvature in some Euclidean or Lorentzian affine 3-space of $\mathbb{L}^4$, again by Remark 8.

Many of the examples with $c = 0$ and non-parallel mean curvature are complete. For instance, if $(g,\omega)$ are the Weierstrass data of a complete minimal surface in $\mathbb{R}^3$ lying in a halfspace with horizontal boundary, then its third coordinate $\text{Re} \int g\omega$ is non-surjective, and thus the above condition C.2 holds. So, $g,\omega$ together with $\varepsilon = -1$ and $a, b \in \mathbb{R}$, $a + \varepsilon b \neq 0$, generate a complete surface of Bryant type in $\mathbb{L}^4$.

As a closing remark, we indicate that if $a = c = 0$ and $b \neq 0$, the resulting class of surfaces admit an integral representation quite analogous to the Weierstrass representation of minimal and maximal surfaces. Indeed, in that situation we have $\vartheta = 0$ in (3.17), so a solution $F : \Sigma \rightarrow \text{SL}(2, \mathbb{C})$ to (3.5) is obtained by substituting $g$ in (4.1) by $\int dg/f$. With this, the system (3.6) can be integrated much in the same way that we did for minimal and maximal surfaces in (4.2). We do not write the final expressions explicitly, as they are straightforward but rather lengthy.

5 Completeness

The well known Calabi-Bernstein theorem [Cal] asserts that the only complete maximal surfaces in $\mathbb{L}^3$ are spacelike planes. Analogously, every complete spacelike CMC-$r$ surface in $\mathbb{S}^3_1(r^2)$ must be a flat totally umbilic example, obtained as the intersection of $\mathbb{S}^3_1(r^2)$ with a degenerate vector hyperplane of $\mathbb{L}^4$ [Aku, Ram]. We remark that both maximal surfaces in $\mathbb{L}^3$ and CMC-$r$ surfaces in $\mathbb{S}^3_1(r^2)$ have non-negative curvature, $K \geq 0$.

We start this section with a simultaneous generalization of these two Bernstein-type theorems:

**Corollary 13** Every complete surface of Bryant type in $\mathbb{L}^4$ with non-negative curvature is a flat surface lying in a degenerate hyperplane of $\mathbb{L}^4$, as described in Remark 6.

**Proof:** Given a surface of Bryant type $\psi : \Sigma \rightarrow \mathbb{L}^4$ with non-negative curvature, we obtain that $\varepsilon = 1$, and so $\langle d\psi, d\psi \rangle = (1 - |g|^2)^2 |\omega|^2 \leq |\omega|^2$. Thus $|\omega|^2$ is a flat metric, which is complete because so is $\psi$. Therefore, the Riemann surface $\Sigma$ must be parabolic, and as $|g| < 1$ we obtain that $g$ is constant. Therefore the Hopf differential $Q$ vanishes identically on $\Sigma$, and the surface must be flat and lie in a degenerate hyperplane of $\mathbb{L}^4$, by Remark 6.

In minimal surface theory, as well as in Bryant surface theory, the study of the complete examples of finite total curvature has been widely developed. Here, we recall...
that a surface $\Sigma$ with non-positive curvature $K \leq 0$ has finite total curvature provided

$$\int_{\Sigma} K \, dA > -\infty,$$

where $dA$ is the area element of the surface.

Our next result shows that, even though there are many complete simply connected surfaces of Bryant type with non-parallel mean curvature, none of them has finite total curvature.

**Theorem 14** Let $\psi : \Sigma \to \mathbb{L}^4$ be a non-flat complete simply-connected surface of Bryant type with finite total curvature. Then $\psi(\Sigma)$ lies in some Euclidean or hyperbolic 3-space of $\mathbb{L}^4$, and its Weierstrass data are given by

$$g(z) = \frac{P_1(z)}{P_2(z)}, \quad \omega = P_2(z)^2 dz, \quad c = a + \varepsilon b = 0,$$

where $P_1(z), P_2(z) : \mathbb{C} \to \mathbb{C}$ are polynomials with no common zeros.

**Proof:** Since $\psi$ is a non-flat immersion then $\varepsilon = -1$, from Corollary 13. On the other hand, using that $\psi$ has finite total curvature, $\Sigma$ must be parabolic, that is, we can assume $\Sigma = \mathbb{C}$, and its Weierstrass data $(g, \omega)$ are meromorphic on $\mathbb{C} \cup \{\infty\}$ (see [Oss]).

As $\omega$ is holomorphic on $\Sigma = \mathbb{C}$, there exist $Q_1(z), Q_2(z), Q_3(z) : \mathbb{C} \to \mathbb{C}$ polynomials, $Q_1(z), Q_2(z)$ without common factors, such that

$$g(z) = \frac{Q_1(z)}{Q_2(z)}, \quad \omega = Q_3(z) dz.$$

Observe that from C.1, the zeros of $g$ of order $k$ correspond to the zeros of order $2k$ of the holomorphic 1-form $g^2 \omega$. So, $g^2 \omega = AQ_1(z)^2 dz$ for a non-zero complex constant $A$. Then, $Q_3(z) = AQ_2(z)^2$ and we can write $g(z) = P_1(z)/P_2(z)$, $\omega = P_2(z)^2 dz$ for $P_1(z) = \sqrt{A}Q_1(z)$, $P_2(z) = \sqrt{A}Q_2(z)$.

Let $R_1, R_2$ be the degrees of $P_1(z), P_2(z)$, respectively. Since $g$ is unique up to isometries of $\mathbb{C} \cup \{\infty\}$, we can suppose that $g(\infty) = \infty$, that is, $R_1 > R_2$. Consequently

$$\deg(\omega) < \deg(g) = \deg(g^2 \omega).$$

Moreover, from (3.3), $f(z)$ is a polynomial such that $\deg(df) = \deg(g^2 \omega)$ if $c \neq 0$ or $\deg(df) = \deg(g \omega)$ if $c = 0$ and $a - b \neq 0$. Hence, if $c \neq 0$ or $a - b \neq 0$ it follows that $\deg(df) \geq \deg(g \omega) = R_1 + R_2$, and so

$$\deg(f) \geq R_1 + R_2 + 1 > R_1 + R_2 - 1 \geq \deg(\omega dg).$$

But the last inequality implies that $\omega dg/f$ is not holomorphic, which is a contradiction. Thus we can conclude that $c = 0$ and $a - b = 0$, which means by Remark 8 that $H$ is parallel, and $\psi$ lies in some Euclidean or hyperbolic 3-space of $\mathbb{L}^4$. This concludes the proof.

$\square$
Appendix

In searching an adequate family which generalizes simultaneously the theories of minimal surfaces in $\mathbb{R}^3$ and Bryant surfaces in $\mathbb{H}^3$, a natural hypothesis is to ask the mean curvature vector to be parallel in the normal bundle, that is, $\nabla^\perp H \equiv 0$. However, this condition is too strong, as it does not generate any new example:

**Proposition 15** Any spacelike surface in $\mathbb{L}^4$ with parallel mean curvature which is not maximal in $\mathbb{L}^4$ must lie in a totally umbilical affine hyperquadric $Q^3$ of $\mathbb{L}^4$.

**Proof:** Let $\psi : \Sigma \to \mathbb{L}^4$ be a non maximal spacelike surface with parallel mean curvature $H$. Then it holds that $\langle H, H \rangle = A$ with $A \in \mathbb{R}$.

If $A = 0$, then $\psi$ is a marginally trapped surface, and we can choose $\nu \in T^\perp(\Sigma)$ such that $\langle \nu, \nu \rangle = 0$ and $\langle \nu, H \rangle = C > 0$. Then it is easy to check that $\{ \eta = (\nu + H)/(2C), \tilde{\eta} = (\nu - H)/(2C) \}$ is an orthonormal frame of $T^\perp(\Sigma)$, and that both $\eta, \tilde{\eta}$ are parallel. This implies that the normal curvature vanishes identically, $R^\perp \equiv 0$. The result follows then by Remark 8.

If $A \neq 0$, let us suppose that $A > 0$ and so $H$ is spacelike (the case $A < 0$ is analogous). Then we can take $\tilde{\eta}$ a normal timelike vector field such that $\{ \eta = H/A, \tilde{\eta} \}$ is an orthonormal frame of $T^\perp(\Sigma)$. As both $\eta, \tilde{\eta}$ are parallel, $R^\perp \equiv 0$. Since the mean curvature vector $H$, given by (2.5), is spacelike, we get $E = 0$ and $E/\lambda = 0$. Now, by the Codazzi and Ricci equations in (2.4) it follows that $p, \tilde{p}$ are holomorphic and $\tilde{p} = kp$, $k \in \mathbb{R}$. The proof of this case finishes following the argument of Remark 8.

The following elementary fact is used in the proof of the representation theorem.

**Lemma 16** Let $f_1, f_2, f_3, f_4 : \Sigma \to \mathbb{C} \cup \{\infty\}$ be meromorphic functions on a Riemann surface, such that $f_1$ and $f_3$ are linearly independent and

$$f_1 \tilde{f}_2 + f_3 \tilde{f}_4 = \tilde{f}_1 f_2 + \tilde{f}_3 f_4.$$  \hspace{1cm} (5.1)

Then there exist constants $a, b \in \mathbb{R}$ and $c \in \mathbb{C}$ such that

$$f_2 = af_1 + cf_3, \quad f_4 = \tilde{c}f_1 + bf_3.$$

**Proof:** Given $z_0 \in \Sigma$ a point where the functions $f_i$ have no poles and $f_3(z_0) \neq 0$, by differentiation of (5.1) we get

$$f_1^{(n)} \tilde{f}_2 + f_3^{(n)} \tilde{f}_4 = \tilde{f}_1 f_2^{(n)} + \tilde{f}_3 f_4^{(n)}$$ \hspace{1cm} (5.2)

at $z_0$, for every $n \in \mathbb{N}$. Thus, all derivatives of $f_4$ at $z_0$ are a linear combination of the derivatives of $f_1, f_2, f_3$ at $z_0$. This shows the existence of $\lambda, \mu, \delta \in \mathbb{C}$ with

$$f_4 = \lambda f_1 + \mu f_2 + \delta f_3.$$  \hspace{1cm} (5.3)
Deriving now (5.2) with respect to $\bar{z}$ at $z_0$ we obtain
\begin{equation}
\frac{\partial}{\partial z} f_1^{(n)} f_2^{(k)} + \frac{\partial}{\partial z} f_3^{(n)} f_4^{(k)} = \frac{\partial}{\partial z} f_1^{(k)} f_2^{(n)} + \frac{\partial}{\partial z} f_3^{(k)} f_4^{(n)}
\end{equation}
for all $k \in \mathbb{N}$. As $f_1, f_3$ are linearly independent, there is some $n_0 \in \mathbb{N}$ such that $f_1^{(n_0)} + \bar{\mu} f_3^{(n_0)} \neq 0$ at $z_0$. Putting this fact together with (5.3) and (5.4) we conclude as before that $f_2$ is a linear combination of $f_1$ and $f_3$. Therefore, there exist complex constants $a, b, c, e \in \mathbb{C}$ such that
\begin{equation}
f_2 = af_1 + cf_3, \quad f_4 = ef_2 + bf_4.
\end{equation}
Now, by (5.1) and (5.5) it holds
\begin{equation}
(a - \bar{a}) |f_1|^2 + (c - \bar{c}) f_1 f_3 + (e - \bar{c}) f_1 \bar{f}_3 + (b - \bar{b}) |f_3|^2 = 0.
\end{equation}
By the linear independence of $f_1$ and $f_3$ this indicates that $a, b \in \mathbb{R}$ and $e = \bar{c}$, what completes the proof.

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