Study of CMOS Sensing System for Radon and Alpha Radiation

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Simple Summary: Radon gas represents one of the main sources of exposure to radioactivity in humans. It’s the second most important cause of lung cancer and the main cause of lung cancer in non-smokers. Commercial radon detectors suffer from low resolution and lack of online monitoring capability, while continuous radon detectors are expensive. This study presents a new semiconductor sensing system for radon and alpha radiation in general, which is fabricated in a standard 180 nm CMOS process with a single polysilicon layer.

Abstract: This study focuses on a CMOS sensing system for Radon and alpha radiation, which is based on a semiconductor device that is integrated monolithically on a single chip with the Readout Circuitry, thus allowing fabrication of a low-power and low-cost sensing system. The new sensor is based on a new mosaic design of an array of Floating Gate non-volatile memory-like transistors, which are implemented in a standard CMOS technology, with a single polysilicon layer. The transistors are electrically combined in parallel and are operated at subthreshold, thus achieving very high sensitivity and reduced noise. The sensing system’s architecture and design is presented, along with key operation concepts, characterization, and analysis results. Alpha and radon exposure results are compared to commercial radon detectors. The new sensor, dubbed TODOS-Radon sensor, measures continuously, is battery operated and insensitive to humidity.

Keywords: radon; alpha radiation; CMOS sensor; silicon radiation detector; ionizing radiation sensor; floating gate transistor

1. Introduction

Radon gas (Rn-222) is a noble, odorless, invisible, and chemically inert radioactive gas that emerges naturally from soils and rocks. It represents one of the main sources of exposure to radioactivity in humans, as it emits alpha particles that may damage lung tissues after inhalation [1]. It is by now well-established that radon gas is the second most important cause of lung cancer and the most important cause of lung cancer in non-smokers, while every 1 in 15 homes in the US suffers from high level of radon concentration [2,3].

The commercial solid state semiconductor radon sensors are reviewed in [4,5] and are implemented in PIN diodes. Low dark current PIN diodes are expensive since large area photodiodes often have defects leading to large dark currents; thus, the sensors are expensive.

The goals of this study were to specify the requirements for Radon gas sensors for smart homes and report how to meet these requirements. It is obvious that for consumer applications, the sensor must be low-cost, small, require low power, and exhibit high sensitivity and reliability. It is obvious to the authors that a standard CMOS technology is mandatory, and a transistor will exhibit higher sensitivity than a PIN diode.

The innovation of the new sensor, dubbed TODOS-Radon Sensor is outlined in the sections below: Section 2.3 presents the new sensor and system architecture and innovations. Device characterization is reported in Section 3.1. Section 3.2 presents experimental results of Radon gas exposure. Section 4 concludes this study.
2. Materials and Methods

2.1. Floating Gate Transistor

The operation principle to sense radiation using Floating Gate (FG) transistor is to first pre-charge the FG using Fowler–Nordheim (FN) Tunneling by applying a large enough electric field on the FG itself [6]. Once the FG is charged, an alpha particle incident will generate e–h pairs that will be separated by an applied electrical field and will discharge the FG [7]. Considering the fact that the threshold voltage of the device is proportional to the charge within the FG as described by Equation (1), an alpha incident will shift the device’s threshold voltage as illustrated in Figure 1.

\[
V_{th} = \phi_{ms} - \frac{Q_{ss,ox}}{C_{ins}} - \frac{Q_{ss,sti}}{C_{sti}} + \frac{Q_{FG}}{C_{sti}} + \frac{Q_{D,max}}{C_{ins}} + 2\phi_f \tag{1}
\]

where \(\phi_{ms}\) is the combined work function difference, \(Q_{ss,ox}\) and \(Q_{ss,sti}\) are the fixed charge densities in the read out oxide and in the control oxide respectively and \(C_{ins}\) is the serial oxide-STI capacitance, given by (see Figure 1a):

\[
C_{ins} = \left( \frac{1}{C_{sti}} + \frac{1}{C_{ox}} \right)^{-1} \tag{2}
\]

\[
\Delta V_{th} = \frac{\Delta Q_{FG}}{C_{sti}} = \frac{N_{e-h} q}{\varepsilon_0 \varepsilon_{ox} A / t_{sti}} \tag{3}
\]

\(N_{e-h}\) is the number of generated e–h pairs by a single alpha incident, \(A\) is the sensing capacitor area and \(t_{sti}\) is the sensing capacitor thickness.

![Figure 1](image_url)

**Figure 1.** FG transistor radiation response illustration: (a) Sensing capacitor charging. (b) e–h generation and separation. (c) Threshold voltage shift: \(C_{sd}\)—sensing area capacitance; \(Q_{inj}\)—injected charge by FN tunneling; \(Q_{ion}\)—ionized charge in an alpha incident.
The number of e–h pairs generated by an alpha incident-$N_{e-h}$, can be calculated by dividing the absorbed energy of a single alpha particle with the e–h pair generation energy. In the proposed Radon Sensor, the STI layer thickness is 3500Å. The absorbed energy is calculated using popular SRIM software [8]. As shown in Figure 2e–h pair generation energy in SiO$_2$ is 17 eV [9].

![SiO$_2$ - 5.5MeV Alpha Absorbed Energy](image)

**Figure 2.** Absorbed energy of a 5.5 MeV alpha particle crossing SiO$_2$ layer with range of 3500Å.

The total absorbed energy within the STI layer is 53.7 KeV, so the number of e–h pairs produced by an alpha particle is:

$$N_{e-h} = \frac{\text{absorbed energy}}{\text{pair generation energy}} = \frac{53.7 \text{ KeV}}{17 \text{ eV}} \approx 3150 \quad (4)$$

### 2.2. C-Sensor

Fabricating an FG transistor in standard CMOS process with a single polysilicon layer was done in the mature, well-proven TowerJazz High Voltage standard 180 nm CMOS Process [10]. TowerJazz had been developing Floating-Gate technology for memory devices, which had been later further developed and extended for direct sensing of radiation and dubbed C-cell. This technology and its reliability have been documented in the open literature [11–16]. The essence of this technology is shown in Figure 3.

![Implementation of C-Sensor in a single-poly CMOS process-schematic cross section](image)

**Figure 3.** Implementation of C-Sensor in a single-poly CMOS process-schematic cross section.

### 2.3. System Implementation

A way to effectively increase the sensing area of the device while keeping high yield, is by using a mosaic array of C-Sensors, where all cells are connected in parallel as shown
in Figure 4. Each cell in the mosaic array is considered as sub-pixel. Although a defect in a single sub-pixel will disable the sub-pixel itself, it will be negligible considering that there are thousands of sub-pixels.

**Figure 4.** Mosaic connectivity: (a) schematics; (b) Layout.

The sensors array is built by 90 × 64 sub-pixels, where each sub-pixel sensing area is 17.5 µm × 12.5 µm, so that the entire array sensing area is 1.12 mm × 1.12 mm.

In order to decrease DC and temperature dependencies, a differential measurement is performed using the concept of active and blind sensor arrays. Considering that alpha particle is a heavy particle, meaning its penetration depth is relatively small, few layers of Kapton film (~50 µm) were used to cover one of the arrays so it could be considered as a blind detector. The device layout is shown in Figure 5.

**Figure 5.** CMOS Sensor device layout.
Another key concept in the proposed architecture is operating in the sub-threshold region, decreasing the power and noise, and also increasing the sensitivity due to the exponential relation \( I = I_{ds} \propto \exp\left(\frac{V_{GS} - V_{th}}{nKT}\right) \) at subthreshold.

\[
i_{sig} = \frac{dI}{dV_{th}} \Delta V_{th} = \frac{q}{nKT} I \Delta V_{th}
\]

Operating in the subthreshold region introduces a linear relation between the output signal and the bias current—\( I \). This point is used as a feature to configure the sensitivity of the device, adapting it to long/short term measurements.

Since the power consumption of the device is dominated by the bias current of the sub-pixels, working in the subthreshold will result in low power consumption which will enable the device to be operated by standard AAA batteries.

An incident alpha particle on the active sensors array will increase current as shown in Equation (5). Subtracting the blind array’s current from the active array’s current provides a differential measurement. Hence, the readout circuit was designed as a precise current subtractor followed by a Trans Impedance Amplifier (TIA), to convert the output signal into output voltage (see Figure 6). The difference between two adjacent samples of the readout voltage, meaning the output voltage derivative, reflects the amount of alpha particles incidents.

\[\Delta V_{th,\text{one}} = \frac{1}{\epsilon_{\text{ox}} \epsilon_{\text{ox}} A / t_{\text{ox}}} = 3150 \cdot 1.6 \cdot 10^{-19} \cdot 8.84 \cdot 3.9 \cdot 10^{-16} \cdot 12.5 \cdot 17.5 \cdot 10^{-16} / 3500 \cdot 10^{-16} = 23\text{mV}
\]

So that the average threshold voltage shift of the entire array is:

\[
\Delta V_{th} = \Delta V_{th,\text{one}} \times \frac{64 \times 90}{3500} = 4\text{\mu V}
\]

3. Results
3.1. Device Characterization

As the first part of the device characterization, the charging mechanism was tested to achieve controllability over the threshold voltage of each sensor array. Figure 7 shows an initial calibration process to shift the sub-pixels threshold voltage into higher and positive values.
Considering Equation (3), the threshold shift of a single sub-pixel caused by a single alpha incident is given in Equation (6).

$$\Delta V_{\text{th},\alpha} = \frac{N_{e-h} q}{\varepsilon_0 \varepsilon_{\text{ox}} A / t_{\text{si}}} = \frac{3150 \cdot 1.6 \cdot 10^{-19}}{8.84 \cdot 3.9 \cdot 10^{-12} / 3500 \cdot 10^{-10}} = 23 \text{ mV}$$ (6)

So that the average threshold voltage shift of the entire array is:

$$\Delta V_{\text{th}} = \frac{\Delta V_{\text{th},\alpha}}{64 \times 90} = 4 \mu\text{V}$$ (7)

3. Results

3.1. Device Characterization

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![Initial Full Arrays IV curve](image1)

![Calibrated Full Arrays IV curve](image2)

Figure 7. Full Arrays IV Curves. (a) initial IV curve; (b) calibrated IV curve.

The need to have such controllability over the threshold voltage comes from the need to: (a) minimize the offset between the arrays, keeping the readout circuit out of saturation, (b) configure high threshold voltage, to increase the e–h separation ratio, and (c) perform a sort of reset stage once the readout circuit is close to saturation.

To characterize the device response to alpha particles radiation, the device was exposed to an alpha source, Po-210, with an alpha emission rate of 925 Bq for two periods of 6 min, with a reset stage before each exposure.

The number of alpha particles reaching the detector is given by:

$$\text{(emission rate)} \cdot \frac{A_D}{2R^2\pi}$$ (8)

where $A_D$ is the detector active sensing area and $R$ is the distance of the detector from the alpha source. Incident alpha ratio is calculated to be $45 \frac{\text{particles}}{\text{sec}}$, so that the expected threshold voltage shift can be calculated by multiplying the threshold voltage shift caused by a single alpha incident, with the number of alpha particles reaching the detector:

$$\Delta V_{\text{th}} = 4 \mu\text{V} \cdot 45 \frac{\text{particles}}{\text{sec}} \cdot \frac{300}{10^{-10}} \approx 52 \text{ mV}$$ (9)
Figure 8 describes the experiment flow and the threshold voltage shift. Table 1 summarizes the results over three different devices, where each device was exposed twice to the alpha source, while a reset stage was performed before each exposure (see Figure 9). Table 1 indicates the bias current of the active device in each exposure, the threshold voltage shift and the readout voltage change through the exposure.

![Initial IV curve](image1)

![IV Curve post first exposure](image2)

![IV Curve After Reset](image3)

![IV Curve post second exposure](image4)

**Figure 8.** Blind and Active arrays IV curve. (a) Initial IV curve. (b) post 6 min exposure IV curve. (c) IV curve after reset. (d) post second 6 min exposure.

**Table 1.** Summary of CMOS Sensor exposure to Alpha radiation.

| #Device | #Measurement | Bias Current | $\Delta V_{th}$ | $\Delta V_{out}$ |
|---------|--------------|--------------|----------------|-----------------|
| 1       | 1            | 103 $\mu A$  | 57 mV          | 1.75 V          |
| 1       | 2            | 102 $\mu A$  | 53 mV          | 1.81 V          |
| 2       | 1            | 90.5 $\mu A$ | 60 mV          | 1.8 V           |
| 2       | 2            | 85 $\mu A$   | 52 mV          | 1.76 V          |
| 3       | 1            | 98 $\mu A$   | 44 mV          | 1.92 V          |
| 3       | 2            | 110 $\mu A$  | 44 mV          | 1.45 V          |
Figure 8. Blind and Active arrays IV curve. (a) Initial IV curve. (b) post 6 min exposure IV curve. (c) IV curve after reset. (d) post second 6 min exposure.

Figure 9. $V_{\text{out}}$ profile during alpha radiation exposure. $V_{\text{out}}$ is measured at the output of the TIA (see Figure 6a).

3.2. Radon Gas Exposure

Four instances of the proposed radon sensors were exposed to radon gas in concentrations of $200 \frac{Bq}{m^3}$ up to $800 \frac{Bq}{m^3}$ for two periods of 9 days and 7 days respectively. The sensor results were compared with commercial detectors: Radon Eye plus and Airthings wave.

Figure 10 presents 27 mV threshold voltage shift after 9 days of radon exposure for a single sensor, while Figure 11 presents the readout circuit derivative of all instances, which represents the voltage difference between two measures caused by alpha particles incident on the active detector. All instances present the same trend, while the distribution of the results reflects each sensor’s sensitivity, which is configured by the sensor’s bias current.

Figure 10. IV curve of blind and active arrays, before and after 9 days of Radon exposure.
Figure 10 presents 27 mV threshold voltage shift after 9 days of radon exposure for a single sensor, while Figure 11 presents the readout circuit derivative of all instances, which represents the voltage difference between two measures caused by alpha particles incident on the active detector. All instances present the same trend, while the distribution of the results reflects each sensor’s sensitivity, which is configured by the sensor’s bias current.

Figure 11. Readout voltage derivative.

Figure 12 shows the comparison between commercial radon sensors to the proposed CMOS sensor. In order to convert the readout voltage into radon concentration, a calibration factor is suggested:

\[
\text{Calibration Factor} = \frac{\text{REF}_{\text{avg}}}{\text{CRS}_{\text{avg}}\cdot I_{\text{bias}}} \tag{10}
\]

Calculating the calibration factor for the compared device in Figure 12, is given by:

\[
\text{Calibration Factor} = \frac{501 \text{ Bq/m}^3}{0.58 \text{ mV} \cdot 96 \mu\text{A}} = 9.09 \text{ Bq/m}^3 \cdot \text{mV/} \mu\text{A} \tag{11}
\]

Figure 12 shows the comparison to commercial radon monitors.

Figure 12. CMOS Sensor comparison to commercial detectors [17,18].
Calculating the calibration factor for the compared device in Figure 12, is given by:

\[
\text{Calibration Factor} = \frac{501 \, \text{Bq/m}^3}{0.58 \, \text{mV}\cdot96 \, \mu A} = 9.09 \, \frac{\text{Bq}}{\text{m}^3}\cdot\text{mV}\cdot\mu A
\] (11)

Figure 13 shows radon measurement of a different CMOS sensor device, with a higher sample rate, at a different period of time. Calibration factor is extracted again:

\[
\text{Calibration Factor} = \frac{180.5 \, \text{Bq/m}^3}{0.22 \, \text{mV}\cdot85 \, \mu A} = 9.6 \, \frac{\text{Bq}}{\text{m}^3}\cdot\text{mV}\cdot\mu A
\] (12)

**Figure 13.** CMOS Sensor comparison to commercial detectors.

### 4. Conclusions

A new sensing system for detecting high energy alpha particle and Radon gas is reported. The system is based on a CMOS IC Sensor which was fabricated in a standard 0.18 μm CMOS process with a single polysilicon layer. The sensor’s unique architecture is based on active and blind mosaic arrays of FG cells, operating in sub threshold region.

The sensing system was tested by exposure to an alpha radiation source in a controlled environment as well as to Radon gas and showed high sensitivity and good correlation to other commercial radon detectors.

In summary, the main advantages of the new Sensing System are: low cost compared to commercial radon sensors, enabling it to be applied in every smart home; low power and may be battery operated using standard AAA batteries; and the integrated digital sensor allows for continuous sampling of radon gas in indoor air providing hourly resolution of Radon levels. The CMOS die is passivated, and therefore insensitive to humidity.

### 5. Patents

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