LOCALIZED STARBURSTS IN DWARF GALAXIES PRODUCED BY
THE IMPACT OF LOW-METALLICITY COSMIC GAS CLOUDS

J. Sánchez Almeida\textsuperscript{1,2}, B. G. Elmegreen\textsuperscript{3}, C. Muñoz-Tuñón\textsuperscript{1,2}, D. M. Elmegreen\textsuperscript{5}, E. Pérez-Montero\textsuperscript{5}, R. Amorín\textsuperscript{6}, M. E. Filho\textsuperscript{1,7,8,9}, Y. Ascasibar\textsuperscript{10}, P. Papaderos\textsuperscript{8,9}, and J. M. Vílchez\textsuperscript{5}

\textsuperscript{1}Instituto Astrofísica de Canarias, E-38200 La Laguna, Tenerife, Spain; jos@iac.es
\textsuperscript{2}Departamento de Astrofísica, Universidad de La Laguna, La Laguna, Tenerife, Spain
\textsuperscript{3}IBM Research Division, T. J. Watson Research Center, Yorktown Heights, NY 10598, USA
\textsuperscript{4}Department of Physics and Astronomy, Vassar College, Poughkeepsie, NY 12604, USA
\textsuperscript{5}Instituto de Astrofísica de Andalucía, CSIC, Granada, Spain
\textsuperscript{6}INAF-Osservatorio Astronomico di Roma, Monte Porzio Catone, Italy
\textsuperscript{7}SIM/CENTRA, Lisbon, Portugal
\textsuperscript{8}Centro de Astrofísica da Universidade do Porto, Porto, Portugal
\textsuperscript{9}Instituto de Astrofísica e Ciências do Espaço, Universidade de Lisboa, Lisboa, Portugal
\textsuperscript{10}Universidad Autonoma de Madrid, Madrid, Spain

Received 2015 July 9; accepted 2015 August 10; published 2015 September 2

ABSTRACT

Models of galaxy formation predict that gas accretion from the cosmic web is a primary driver of star formation over cosmic history. Except in very dense environments where galaxy mergers are also important, model galaxies feed from cold streams of gas from the web that penetrate their dark matter halos. Although these predictions are unambiguous, the observational support has been indirect so far. Here, we report spectroscopic evidence for this process in extremely metal-poor galaxies (XMPs) of the local universe, taking the form of localized starbursts associated with gas having low metallicity. Detailed abundance analyses based on Gran Telescopio Canarias optical spectra of 10 XMPs show that the galaxy hosts have metallicities around 60% solar, on average, while the large star-forming regions that dominate their integrated light have low metallicities of some 6% solar. Because gas mixes azimuthally in a rotation timescale (a few hundred Myr), the observed metallicity inhomogeneities are only possible if the metal-poor gas fell onto the disk recently. We analyze several possibilities for the origin of the metal-poor gas, favoring the metal-poor gas infall predicted by numerical models. If this interpretation is correct, XMPs trace the cosmic web gas in their surroundings, making them probes to examine its properties.

Key words: galaxies: abundances – galaxies: dwarf – galaxies: evolution – galaxies: formation – galaxies: structure – intergalactic medium

1. MOTIVATION

Metals are primarily produced by stars. Thus, metal-poor objects are chemically unevolved, providing a gateway to sample physical conditions and processes characteristic of early phases in the universe. XMP galaxies, defined as having less than one-tenth of the metals in the solar composition (Kunth & Östlin 2000), are rare in the local universe (<0.2% of all galaxies; Morales-Luis et al. 2011). They tend to be dwarf galaxies, but unlike other dwarfs (Kobulnicky & Skillman 1996), their chemical composition may not be uniform (Papaderos et al. 2006; Izotov et al. 2009; Levesque et al. 2011; Sánchez Almeida et al. 2013, 2014b), with the low metallicity only in regions of intense star formation. Such non-uniformity is particularly revealing because the timescale for azimuthal mixing in normal galaxies is short, of the order of a fraction of the rotational period (e.g., de Avillez & Mac Low 2002; Yang & Krumholz 2012). It would imply that the metal-poor gas in XMPs was recently accreted from a nearly pristine cloud, providing long-sought evidence for cosmic accretion that has been difficult to observe directly (Sánchez Almeida et al. 2014a). For example, gas flows have been observed on the periphery of distant galaxies using absorption lines in the spectra of background sources. However, gas flows are not significantly lower in metals than the galaxy itself, suggesting gas recycled from previous outflows, and therefore dynamically detached from the large-scale cosmic web (Churchill et al. 2013; Lehner et al. 2013; Liang & Chen 2014).

Here, we put on firm observational bases the existence of chemical inhomogeneities in most XMPs, thus unequivocally showing their star formation to be feeding from external metal-poor gas. Measuring chemical inhomogeneities in XMPs is technically challenging, requiring spectra of high quality from faint regions. We used the 10.4 m Gran Telescopio Canarias (GTC)\textsuperscript{11} (Section 2) to measure oxygen abundance along the major axes of 10 XMP galaxies (Table 1) with a robust variant of the direct method (Pérez-Montero 2014; see Section 3). In 9 out of 10 cases, sharp metallicity drops were found (Section 4). The origin of the metal-poor gas is analyzed in Section 5, leaving the gas infall from the web predicted by numerical models as the most convincing hypothesis.

2. OBSERVATIONS

The 10 galaxies in Table 1 represent 7% of all known XMPs when the project began (Morales-Luis et al. 2011), and they were chosen to be representative of the XMP family. They were observed with the long-slit mode of the spectrograph OSIRIS at GTC. We used an intermediate-resolution grism that renders the spectral range needed for abundance analysis in a single exposure (from 3700 to 7000 Å) with a dispersion of 7.8 Å pixel\textsuperscript{−1}, enough to spectrally resolve the required emission lines. The width of the spectrograph slit was set to 1 arcsec, which matches the typical seeing during

\textsuperscript{11} http://www.gtc.iac.es/GTChome.php
Table 1
Extremely Metal Poor Galaxies Selected in This Study

| Name           | 12 + log(O/H) | log \(M_\odot\) | \(D^d\) | H \(\alpha\) Size | \(\Delta Z^f\) |
|----------------|---------------|-----------------|----------|------------------|----------------|
| J020549.3−094920.3 | 7.58 ± 0.14   | 8.26 ± 0.48     | 26.4 ± 1.8 | 0.18 ± 0.03      | Y              |
| J030331.3−010947.1 | 7.68 ± 0.13   | 8.33 ± 0.33     | 125 ± 9   | 0.72 ± 0.12      | Y              |
| J031300.0−000612.2 | 7.58 ± 0.17   | 7.67 ± 0.31     | 119 ± 9   | 0.35 ± 0.12      | Y              |
| J082555.0−353231.0 | 7.49 ± 0.11   | 6.04 ± 0.03     | 9.6 ± 0.7 | 0.076 ± 0.009    | Y              |
| J094416.6+541134.4 | 7.43 ± 0.15   | 7.05 ± 0.05     | 23.1 ± 1.6 | 0.16 ± 0.2       | Y              |
| J113202.4+572245.2 | 7.59 ± 0.15   | 7.53 ± 0.48     | 22.5 ± 1.6 | 0.093 ± 0.021    | Y              |
| J144506.3+501802.0 | 7.70 ± 0.10   | 6.71 ± 0.07     | 23.6 ± 1.7 | 0.31 ± 0.02      | N              |
| J115133.3−022221.9 | 7.32 ± 0.15   | 6.17 ± 1.0      | 12.6 ± 0.9 | 0.12 ± 0.01      | Y              |
| J210455.0−003522.0 | 7.33 ± 0.10   | 6.19 ± 0.05     | 21.8 ± 1.6 | 0.094 ± 0.021    | Y              |
| J230210.0+004939.0 | 7.54 ± 0.12   | 8.39 ± 0.33     | 138 ± 10  | 0.47 ± 0.13      | Y              |

Notes.

\(^a\) R.A. and decl. in J2000 coordinates.

\(^b\) Oxygen abundance at the starburst.

\(^c\) Stellar mass of the galaxy from MPA-JHU given in units of \(M_\odot\).

\(^d\) Redshift-based galactocentric distance, in Mpc, from NED (NASA/IPAC Extragalactic Database).

\(^e\) Seeing-corrected FWHM size of the starburst, in kpc, as inferred from the H\(\alpha\) emission.

\(^f\) Does the galaxy show chemical inhomogeneities? Yes or No. Note that the only galaxy without chemical inhomogeneities is also the object of the largest metallicity.

3. DETERMINATION OF METALLICITIES, STELLAR MASSES, AND STAR FORMATION RATES (SFRs)

H\(\alpha\) regions are photoionized by young stars. Their emission-line spectrum allows the determination of the oxygen abundance independently of the details of the ionization source and the geometry of the nebula (Osterbrock 1974; Stasińska et al. 2012). The spectrum contains a number of collisionally excited emission lines that empirically provide electron density and electron temperature. Given temperature and density, atomic properties yield the expected emission per ion; therefore, the total flux of a particular emission line directly gives the total number of ions in the nebula responsible for such emission. The case of oxygen is particularly favorable since, in H\(\alpha\) regions, oxygen is in the form of O\(^+\) and O\(^{++}\), and both species produce emission lines. The total oxygen abundance results from adding the contribution of O\(^+\) and O\(^{++}\). A drawback of this technique (the so-called direct method, DM) is the need to measure fluxes of weak lines often at the noise level. Unfortunately, this problem affects the faint tails of our XMPs (Figures 1 and 2). We circumvent the difficulty using HII-CHI-mistry\(^{13}\) (Pérez-Montero 2014), a procedure that compares the brightest optical emission lines with predictions of a grid of photoionization models covering a wide range of possible physical conditions. The resulting metallicity is calculated as the \(\chi^2\)-weighted mean of the metallicities in the models, with the \(\chi^2\)-weighted standard deviation providing the error. This method leads to an oxygen abundance that is consistent with the DM determination within 0.07 dex, i.e., within the intrinsic uncertainty of the DM (Pérez-Montero 2014). The agreement implies that HII-CHI-mistry inherits the good properties of DM so that the geometry and other assumptions needed to build the grid of models do not influence the inferred oxygen abundance. Provided that the emission lines are produced by photoionization, the agreement between HII-CHI-mistry and DM guarantees that irrespectively of the source of ionizing photons, our abundance determinations along the XMP major axes are reliable and can be inter-compared. Photoionization is a good approximation even for the faintest parts of XMPs, as attested by the weakness of the line [N\textsc{ii}] \(\lambda6583\) Å. In the cases where we have been able to measure it, its flux is less than one-tenth of the H\(\alpha\) flux, which, according to the BPT diagnostic diagram (Baldwin et al. 1981), discards ionization by active galactic nuclei, shocks (Kewley et al. 2013), and evolved stars (Cid Fernandes et al. 2011).

The stellar masses of the galaxies given in Table 1 were determined by the MPA-JHU collaboration from galaxy-integrated magnitudes (Brinchmann et al. 2004; Salim et al. 2007). Broadband magnitudes from the SDSS-DR7 database were compared with a grid of theoretical galaxy spectra (Bruzual & Charlot 2003) spanning a large range in star formation histories. From the difference between observed and theoretical magnitudes, a likelihood distribution for the mass of each galaxy is estimated. The median and the dispersion of this distribution are used for the stellar mass and its error.

SFRs are needed for the discussion of the results (e.g., Figures 1 and 2). We inferred them from the observed H\(\alpha\) flux as described by Kennicutt (1998). The underlying hypothesis is that H\(\alpha\) quantifies the number of ionizing photons produced by young (<20 Myr) massive (>10 \(M_\odot\)) stars, and so its flux scales with the current SFR.

4. RESULTS

HII-CHI-mistry provides the abundance of oxygen along the major axis of the galaxies. In 9 out of the 10 cases, sharp metallicity drops were found (examples are given in Figures 1 and 2, with the only exception having constant metallicity pointed out in Table 1). Remarkably, the metallicity drops occur coinciding with starburst regions (the solid lines in Figures 1 and 2, bottom panels).

A summary plot with the metallicity of the starbursts and the underlying galaxies is given in Figure 3. The abundance of the starburst was determined as the lowest abundance, whereas the abundance of the host was computed by averaging all the values in a region of 2.5 arcsec around the position with the
The host galaxies have metallicities higher than the starbursts by factors of 3–10 where their surface SFRs are smaller by factors of 10–100. According to the empirically established scaling between surface SFR and gas surface density (the so-called Kennicutt–Schmidt law; e.g., Kennicutt 1998; Yang & Krumholz 2012; Cormier et al. 2014), the observed starbursts imply large gas reservoirs with surface densities between $10^2$ and $10^3 \, \text{M}_\odot \text{pc}^{-2}$ (Figure 3, upper axis).

According to simple chemical evolution models, the starbursts can evolve to become the host only if they expel 80% of the gas and metals to the intergalactic medium (blue solid line in Figure 3). If the evolution proceeds as a closed box, and it begins from pure metal-free gas, the fraction of mass in metals of the interstellar medium $Z$ depends only on the gas that remains (e.g., Edmunds 1990), so that if a region of metallicity $Z_1$ and gas mass $\Sigma_{\text{gas}1}$ evolves to reach a metallicity $Z_2$ of gas mass $\Sigma_{\text{gas}2}$, then $Z_1 - Z_2 = Y \ln(\Sigma_{\text{gas}2}/\Sigma_{\text{gas}1})$ with $Y$ the stellar yield, i.e., the mass of new metals eventually ejected...
per unit mass locked into stars. The red solid line in Figure 3 shows the evolution expected for a region at the red bullet assuming a typical Oxygen yield (0.004; Meynet & Maeder 2002). It renders metallicities too high to explain the metallicity of the host galaxy by closed-box evolution of the starbursts. If gas is allowed to leak out of the system, the relationship between gas mass and metallicity remains formally as for the closed-box model, replacing the stellar yield with an effective yield that involves the fraction of mass in stars that returns to the interstellar medium, as well as the so-called mass loading factor, which gives the outflow rate in terms of the SFR (e.g., Sánchez Almeida et al. 2014a): \[ Y_{\text{eff}} = Y (1 - R)/(1 - R + w). \] The blue solid line in Figure 3 corresponds to \( Y_{\text{eff}} = 0.2 \) Y, or to \( w = 3.2 \) given \( R = 0.2 \), which implies that for each solar mass of gas locked into stars, four solar masses of gas flow out of the system.

Figure 4 shows the scatter plot of metallicity versus galaxy stellar mass for both the host and the starburst. Our objects are outliers of the mass–metallicity relationship (MZR) observed in the local universe. The green and blue lines in Figure 4 represent MZRs for the range of low-mass galaxies (Lee et al. 2006) and high-mass galaxies (Tremonti et al. 2004), respectively, with the shaded areas portraying the observed standard deviation at a given mass. The high-mass metallicities have been decreased by 0.2 dex to make the strong-line method of the original work consistent with the DM scale used here (Kewley & Ellison 2008). The MZR is based on the metallicity of the brightest parts of the galaxies; therefore, XMPs must be represented by the metallicity of their starbursts, which clearly lie below the observed MZR (Figure 4). The scatter at a given mass is known to be related to the SFR (Lara-López et al. 2010; Mannucci et al. 2010), so that the higher the SFR the lower the metallicity. The pink area in Figure 4 represents the spread in metallicity to be expected if the SFR changes by \( \pm 10 \) times the mean SFR given by the red solid line (Andrews & Martini 2013), and this dependence naturally accounts for the large difference in metallicity between starburst and host. XMPs are going through an active star formation phase, with starbursts of low metallicity. The end of this phase will disclose a galaxy having properties close to those of the host, i.e., a stellar mass similar to the present XMP but with low SFR and high metallicity.

5. DISCUSSION

The XMPs studied here have an off-center starburst of low metallicity; otherwise, they are rather normal disk-like galaxies with exponential light profiles, modest rotational velocities, and significant random motions (Elmegreen et al. 2012; Sánchez Almeida et al. 2013). Under these conditions, the timescale for azimuthal gas-mixing is rather short (see Section 1); therefore, the metal-poor gas associated with the starburst must have arrived to the galaxy recently.

The mass surface density at the starburst is unusually large for the outer parts of dwarf irregular galaxies (Cormier et al. 2014) by a factor of \( \sim 10 \). If the excess mass is from a recently accreted cosmic cloud, then it mixed with the ambient gas and diluted the ambient metallicity by the same factor, \( 10:1 \). This dilution accounts for the metallicity drop in the starbursts. The accretion time is constrained by the localization of the star formation, which is not sheared or spiral-like. The confinement of the enhanced star formation to local regions several hundred parsecs in diameter (Table 1) implies that the accretion time is shorter than the orbital time. For small galaxies like these, the orbital time over one radius is \( \sim 100 \) Myr. Accretion lasting for \( \sim 100 \) Myr or less corresponds to an impacting cloud with a vertical extent that is similar to the radius of the galaxy or several kiloparsecs, considering that the infall speed is likely comparable to the rotation speed since both are set by the gravitational potential. From these considerations, we conclude the following scenario for XMP galaxies with locally low metallicities studied here. Sometime within a period of \( \sim 100 \) Myr, a cosmic cloud with low metallicity impacted the outer region of a regular dwarf galaxy disk, compressing it, diluting the original metals, and triggering excess star formation, all by factors of 10. This picture is consistent with the current ideas on how dwarf galaxies grow through gas accretion (Section 1 and Finlator & Davé 2008; Dekel et al. 2009; Schaye et al. 2010), even though the interaction between gas streams and forming disks remains to be properly understood (Ceverino et al. 2010; Verbeke et al. 2014).

Explanations alternative to the cosmic origin of the gas cannot be completely ruled out, although they are disfavored by observations. Major mergers are not obvious in the optical images, and can be further discarded, arguing that XMPs reside in low-density environments, with their large \( \text{H}_1 \) envelopes showing only moderate distortions (Filho et al. 2015). Minor gas-rich mergers could explain localized metallicity drops, provided that the gas in the accreted galaxy has a metallicity lower than the metallicity at the drop. A minor-merger scenario differs from a gas accretion scenario only if the merging galaxy has enough stars and dark matter. The additional mass would generally be much larger than the merging gas mass alone, and the distortion on the host would be larger as well. Merging stars should also penetrate the host disk, unlike the merging gas, and these stars would orbit differently than the host stars, possibly producing star streams in the host halo. Since neither large \( \text{H}_1 \) distortions nor stellar streams are evident in XMPs (see, e.g., Figures 1 and 2), the hypothetical merging galaxy must have little dark matter and stars and, consequently, it could be likened to a gas stream. On the other hand, model cosmic gas
densest cores. Streams are often clumpy, with some stars forming in the densest cores (e.g., Dekel et al. 2009), in which case the accretion of a clumpy stream is equivalent to a merger with a small very gas-rich, dark-matter-poor satellite. The difficulty in distinguishing a pure gas accretion event from a gas-rich minor merger can be linked to the problem of finding an operational definition of galaxy at the low-mass end of galaxy masses (see Forbes & Kroupa 2011; Willman & Strader 2012). Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.

In addition, the accreted material is expected to spread out over their immediate surroundings, providing a new method to distinguishing a pure gas accretion event from a gas-rich minor merger. Yet, another possibility for feeding XMPs could be gas accretion from the galaxy halo induced by galactic fountains (Fraternali & Binney 2008; Fraternali et al. 2014). The metal-rich gas ejected from the stellar disk acts as a catalyst when it becomes mixed with the hot metal-poor halo gas. The mixing catastrophically decreases the halo cooling time, and mass from the halo is dragged along when the fountain gas falls back onto the disk. Several observations disfavor this possibility, though. The mass dragged along is only a fraction of the original mass, and so the gas falling in is not very metal poor.