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Extended SO\textsubscript{2} outgassing from the 2014-2015 Holuhraun lava flow field, Iceland

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Abstract

The 2014-2015 Holuhraun eruption was the largest fissure eruption in Iceland in the last 200 years. This flood basalt eruption produced ~1.6 km$^3$ of lava, forming a lava flow field covering an area of ~84 km$^2$. Over the six-month course of the eruption ~11 Mt of SO$_2$ were released from the eruptive vents as well as from the cooling lava flow field. This work examines the post-eruption SO$_2$ flux emitted by the Holuhraun lava flow field, providing the first study of the extent and relative importance of the outgassing of a lava flow field after emplacement.

We use data from a scanning differential optical absorption spectroscopy (DOAS) instrument installed at the eruption site to monitor the flux of SO$_2$. In this study, we propose a new method to estimate the SO$_2$ emissions from the lava flow field, based on the characteristic shape of the scanned column density distribution of a homogenous source close to the ground. Post-eruption outgassing of the lava flow field continued for at least three months after the end of the eruption, with SO$_2$ flux between <1 kg/s and 9 kg/s. The lava flow field post-eruption emissions were not a significant contributor to the total SO$_2$ released during the eruption, however the lava flow field was still an important polluter and caused high concentrations of SO$_2$ at ground level after lava effusion ceased.

Keywords SO$_2$, lava, DOAS, post-eruption outgassing
Introduction

A dyke-fed basaltic fissure eruption from the Bárðarbunga volcanic system occurred from 31 August 2014 until 27 February 2015 (Icelandic Meteorological Office (IMO), 2015a). The fissure was located to the north of the Vatnajökull icecap, forming the 2014-2015 Holuhraun lava flow field (Figure 1). The eruption produced $1.6 \pm 0.3$ km$^3$ of lava, forming an $84.1 \pm 0.6$ km$^2$ lava flow field (Gíslason et al., 2015). This classifies the eruption as a flood basalt eruption following Thordarson and Larsen (2007). This makes the Holuhraun eruption the most voluminous effusive eruption in Iceland since the 1783-1784 Laki eruption (Schmidt et al., 2015).

This work provides a new approach for measuring SO$_2$ emissions from a lava flow field after emplacement. Several studies have investigated the impacts of the Holuhraun eruption on populations and the environment (e.g. Gislason et al., 2015; Ilyinskaya et al., 2017). Establishing the potential duration and magnitude of post-eruptive SO$_2$ outgassing will increase resilience in vulnerable communities related to the health hazards caused by SO$_2$ after future eruptions. We additionally discuss the physical processes of cooling and fracturing (e.g. Keszthelyi and Denlinger, 1996; Kattenhorn and Schaefer, 2008; Patrick et al., 2004; Wittmann et al., 2017) that contribute to post-emplacement lava outgassing.

Volatile outgassing during and after eruption

Gases dissolved within magma are transferred to the atmosphere by degassing. The term “degassing” is used to describe the process by which a magma loses its volatiles (Burgisser and Degruyter, 2015). This includes exsolution of gas from the melt, gas segregation, and outgassing (e.g. Bottinga and Javoy, 1991; Sparks, 2003; Palma et al., 2008). Following
Burgisser and Degruyter (2015) and Palma et al. (2008; 2011) we define “outgassing” only as the release of this gas to the atmosphere.

Basaltic fissure eruptions release large volumes of SO$_2$ into the atmosphere. Basaltic magmas have a high sulphur yield, which is typically two to four times higher than silicic magmas (Thordarson et al., 2003). As a result, SO$_2$ is released during basaltic flood eruptions by a two-stage degassing process: from the magma as it rises through the conduit and erupts at the vent, and from lava flows during and after their emplacement (Walker, 1989; Thordarson et al., 1996; Thordarson et al., 2003). Gases released at the vent, and associated with degassing of volatiles as the magma approaches the surface, contribute to an eruption plume. Gases released from lava flows will remain close to ground level forming a low-level haze (Figure 2).

Swanson and Fabbi (1973) studied volatile loss during isothermal flowage of pahoehoe lava in lava tubes at Mauna Ulu, Hawaii. They observed that the majority of sulphur was released soon after eruption (60% was lost during flowage over a distance of 12 km), but lava flows were observed to continue outgassing for at least two to four hours after solidification (Swanson and Fabbi, 1973). In contrast, the Holuhraun lava flow field exhibits a wide range of lava morphologies, varying from pāhoehoe to ‘a’ā (Pedersen et al., 2017) and, as demonstrated by this work, continued to release sulphur volatiles for several months after emplacement.

In another study on Hawaiian flows, Bottinga and Javoy (1991) proposed that lava flows degas volatiles during transportation because their temperature decreases. The resulting supersaturation of volatiles within the lava increases, causing exsolution resulting in bubble formation. Cashman et al. (1994) proposed that volatiles were primarily released from active lava flows by the rise and escape of bubbles already contained within the lava. Sparks and Pinkerton (1978) noted that as lava is degassed, exsolution of gas results in undercooling of the lava. This leads to crystallisation, and causes an increase in viscosity and yield strength (Sparks
and Pinkerton, 1978). Mechanical processes during the flowage of lava then cause the solidified crust to fracture (e.g. Polacci and Papale, 1997; Soule and Cashman, 2004), allowing further gas to be released.

Many studies have investigated the degassing of lava flows as they are actively flowing from the vent, however little discussion exists about the processes of extended outgassing following lava flow emplacement. The studies of Bottinga and Javoy (1991) and Cashman et al. (1994) both examine the release of gases from flowing lava and are therefore insufficient to explain the continued flux of SO$_2$ from the Holuhraun lava flow field for several months after its emplacement and solidification. We thus, here, consider the physical processes which occur over the longer time period that was captured by our outgassing measurements.

**Volcanic SO$_2$ monitoring**

SO$_2$ emission measurements have been fundamental to volcano monitoring since the development of the correlation spectrometer (COSPEC) in the 1970s (Moffat and Millan, 1971; Stoiber et al., 1983). Malinconico (1979) observed that increases in SO$_2$ flux at Mt Etna, Italy, corresponded to increases in volcanic activity, suggesting that SO$_2$ fluctuations can be used to predict eruptions. SO$_2$ flux has since been observed to correlate with volcanic activity (Soufrière Hills, Montserrat; Edmonds et al., 2003), magma extrusion rate (Unzen, Japan; Hirabayashi et al., 1995), seismicity during explosive activity (Merapi, Indonesia; Jousset et al., 2013) and lava lake activity (Villarrica, Chile; Palma et al., 2008, and Erebus, Antarctica; Kyle et al., 1994).

Differential optical absorption spectroscopy (DOAS) was also developed in the 1970s, as a method for measuring atmospheric gases (Platt et al., 1979). It has since been used to measure
SO₂ emissions of volcanoes worldwide (e.g. Edner et al., 1994; Weibring et al., 1998; McGonigle et al., 2002), and has now superseded COSPEC as the primary volcanic SO₂ flux monitoring technique (Galle et al., 2002; Bobrowski et al., 2010). Measurement of volcanic plumes was further enhanced by the development of the SO₂ camera (Mori and Burton, 2006; Bluth et al., 2007), providing high temporal resolution SO₂ flux measurements and allowing SO₂ heterogeneity within the plume to be quantified (Bluth et al., 2007). SO₂ measurement techniques have been employed to determine SO₂ emissions within volcanic plumes, both during eruptive episodes (e.g. Jousset et al., 2013; Gíslason et al., 2015) and during passive degassing (e.g. McGonigle et al., 2002; Sawyer et al., 2008), however this work provides the first attempt at using DOAS to measure the SO₂ released by an emplaced lava flow field.

The 2014-2015 Holuhraun eruption

Over the six month course of the Holuhraun eruption, 11 ± 5 Mt of SO₂ were emitted (Gíslason et al., 2015), with average emission rates of 400 kg/s, and peaks of over 1000 kg/s (Barsotti et al., 2015; Gauthier et al., 2016). The majority of gases were released from the eruption fissure and contributed to an eruption plume that contained very little ash but was concentrated in SO₂ and H₂O (Gíslason et al., 2015). In addition to the principal eruption plume, SO₂ was also released directly from the lava flows, forming a low-level haze of SO₂. A surveillance flight on 4 November 2014, for example, revealed a distinct two-layered gas cloud that included a low-level “haze” of H₂O (both magmatic and meteoric), SO₂ and other volcanic gases. This haze rose from the lava flow field to an elevation of 700 m above ground level, with the eruption plume ascending to an elevation of 1500-2500 m (Figure 3; IMO, 2014). High concentrations of SO₂ were measured at ground level throughout Iceland, and in many communities the health standard of 350 µg/m³/hr was exceeded, posing health risks to the
population (Gíslason et al., 2015, Ilyinskaya et al., 2017). After the eruption ceased at the end of February 2015, high concentrations of SO$_2$ continued to be detected at ground level near to the eruption site, but dropped to background levels at communities downwind (IMO, 2015a; Umhverfisstofnun, 2016). The aim of this work was to measure the SO$_2$ released from the cooling lava flow field after the eruption, and to examine the processes that facilitate an emplaced, cooling lava flow field to outgas over a prolonged period.

**Methods**

The methodology used to derive the SO$_2$ flux from the lava flow field was based on the same principle as is used in COSPEC and MobileDOAS measurements of an elevated gas plume (e.g. Stoiber et al., 1983; Galle et al., 2002). Following this approach, the vertical column density (VCD) of a gas is measured using absorption spectroscopy with the sky as the light source (Platt and Stutz, 2008). By traversing under the plume in a direction approximately perpendicular to the plume propagation, and integrating the obtained vertical column amounts, the total number of molecules in a cross-section of the plume may be determined. After multiplication with the plume speed the total gas emission is calculated (e.g. Edner et al., 1994; Galle et al., 2002; McGonigle et al., 2002; Edmonds et al., 2003; Platt et al., 2015).

In our application, a stationary scanning DOAS instrument (ScanDOAS; Galle et al., 2010) was used to determine the VCD of the low-level haze of gas emitted by the lava flow field. To make this measurement we assume that the lava flow field produces a gas layer close to the ground with uniform thickness. We also assume that this layer has a width that is equal to the maximum width of the lava flow field, measured in a direction perpendicular to the wind direction, the “effective haze width”. Using this geometry, with a ScanDOAS instrument we can determine the VCD at the instrument location, and using the known wind direction and
shape of the lava flow field we can derive the effective haze width. The gas flux can then be
determined by multiplying the VCD at the instrument site with the effective haze width and
with the wind speed.

Measurement geometries

A ScanDOAS scans the sky from horizon to horizon in a vertical plane approximately
perpendicular to the plume propagation, recording radiance spectra of the diffused UV solar
radiation received at each angle (Edmonds et al., 2003). The “slant column density” (SCD) of
SO$_2$ from each spectrum is calculated by differential optical absorption spectroscopy (DOAS),
and the integral of column densities at all angles is then multiplied by the plume speed to obtain
the flux of SO$_2$ (e.g. Stoiber et al., 1983; Galle et al., 2002; McGonigle et al., 2002). The
instrument typically makes one scan approximately every five minutes, each scan being
composed of 26 spectra.

In a variation on this “flat” scan geometry, the scan is made over a conical surface with its tip
at the instrument and its base through the gas source. The main advantage of this “conical”
geometry is that a wider range of plume directions may be covered by a single instrument
(Galle et al., 2010). In this study, a conical geometry with an opening angle of 60° was used.
The instrument used in our study is a modification of the standard NOVAC-Mark I instrument
(NOVAC: Network for Observation of Volcanic and Atmospheric Change; Galle et al., 2010),
with a non-rotating, cylindrical external hood made of quartz, and a UV sensitive OceanOptics
MAYA Pro spectrometer (http://oceanoptics.com/wp-content/uploads/OEM-Data-Sheet-
Maya2000Prov3.pdf).
SO$_2$ emissions from a vent will contribute to an elevated eruption plume, which will be identified on a DOAS scan as a concentrated distribution of SO$_2$ column densities, with higher columns observed at elevation angles in which the scanner detects the bulk of the plume (Figure 4a). In contrast, SO$_2$ emissions from a lava flow field or a grounded plume will form a dispersed low-level haze. This will have a characteristic trough shape on a DOAS scan (Figure 4b). When scanning through a low-level haze of SO$_2$, the optical path will be greater at low elevation angles close to the ground, producing a greater SO$_2$ slant column density. The minimum slant column densities of SO$_2$ will be observed at the highest elevation angles, where the shortest path through the haze is sampled (Figure 4b). With flat geometry this occurs at zenith, but for the conical geometry used here this condition is met at 30° from zenith.

**Determination of the SO$_2$ vertical column density at the measurement site**

If we assume a layer of gas close to the ground, and flat scan geometry, then the slant column densities (SCD) can be expressed as:

$$SCD = \frac{VCD}{\cos(\alpha)}$$  \hspace{1cm} (1)$$

where VCD is the vertical column density through the sampled layer and $\alpha$ is the scan angle measured from the zenith.

In a standard evaluation of a DOAS spectrum, the spectrum measured through the gas plume is divided by a clean air reference spectrum, typically obtained in a direction with no gas (Galle et al., 2002; Edmonds et al., 2003). In this way spectral features related to the sky spectra as well as instrumental features are cancelled out and an absorption spectrum of the gas plume is obtained. When scanning through a low-level haze, all spectra, including the reference spectrum, will contain SO$_2$. Our goal was therefore to determine the VCD without having
access to a clean air reference spectrum. In this case, the derived slant column was the
difference between the slant column of the measured spectrum and the slant column of the
reference spectrum by which it was divided. For a flat geometry this can be achieved by
dividing a spectrum taken at $\alpha = 60^\circ$ with a reference spectrum taken at $\alpha = 0^\circ$. The resulting

$$\text{SCD}_{\text{diff}} = \text{VCD}/0.5 - \text{VCD}/1 = \text{VCD},$$

provides the required VCD. For a conical geometry the same method can be applied. However here the relation between
VCD and SCD becomes:

$$SCD = \frac{VCD}{[\cos(\alpha) \times \sin(\beta)]}$$

(2)

With the conical angle $\beta = 60^\circ$, the difference between the SCD taken at angles $\alpha = 60^\circ$ and $\alpha = 0^\circ$ becomes:

$$\text{SCD}_{\text{diff}} = \frac{VCD}{0.866}$$

(3)

Therefore, to obtain the VCD at the instrument location we multiply the average SCD obtained
at angles $\alpha = 60^\circ$ and $\alpha = 0^\circ$ by 0.866. The SCD and VCD are, by convention, expressed in
ppm*m (parts per million-metre). To calculate SO$_2$ emission, these are converted to kg/m$^2$ by
applying the ideal gas law and multiplying the VCD (in ppm*m) by $2.66 \times 10^{-6}$.

Field configuration and SO$_2$ flux calculation

ScanDOAS instruments were installed at the Holuhraun eruption site by the Icelandic
Meteorological Office and Chalmers University to measure SO$_2$ emissions during the eruption.
Instrument MAYP111126 (cross symbol in Figure 1) was installed at its location to the east of
the lava flow field on 4 March 2015, where it was oriented towards the west and optimally
located for viewing of the lava flow field. This occurred after the eruption ended on 27 February,
and data from this spectrometer were used to calculate the post-eruptive SO$_2$ flux from the lava
flow field.

All scans from instrument MAYP111126, from its installation on 4 March 2015 until 31 May
2015, were examined to identify those with the characteristic trough shape that represented the
anticipated view through a low-level haze of SO$_2$ from the lava flow field (Figure 4c). Each
“haze” scan was first re-evaluated between 310 nm and 325 nm using the NovacProgram
software (version 1.82; Galle et al., 2010), including additional spectra in the DOAS fitting
process to account for the presence of ozone in the atmosphere and the “Ring effect” (i.e. the
“filling-in” of deep absorption features, such as the solar Fraunhofer lines, in a measured
atmospheric spectrum, which is caused by inelastic (Raman) scattering of light by air molecules
(Grainger and Ring, 1962)). After calculating the column densities, a quality check was applied
to ensure that scans retained the characteristic trough shape of the anticipated lava flow field
“haze” scans. To ensure a symmetrical trough shape, scans were excluded if slant column
densities at +60° and -60° differed by more than 20%. To calculate the SO$_2$ flux (in kg/s)
recorded by each scan, the measured VCD of SO$_2$ (in kg/m$^2$) was multiplied by the effective
width of the SO$_2$ haze (in m) and the haze speed (in m/s), assumed to be equivalent to the wind
speed measured at a co-located meteorological station.

As mentioned above, the effective width of the low-level haze produced by the lava flow
field was assumed to be equal to the total width of the lava flow field, measured in a direction
perpendicular to the wind direction (Figure 5). Because the lava flow field has an irregular
shape, this width varies with different wind directions, and can be measured from a map
showing the extension of the lava flow field. The minimum effective haze width occurs when
the wind blows along the longer dimension of the lava flow field, and as wind direction
deviates from this the effective haze width increases (Figure 5).
Wind speed and direction were recorded every 10 minutes using a meteorological station installed at the eruption site. Scans were filtered by wind direction, with scans included only when the wind was blowing towards the instrument (± 20°).

Lava flow field radiant heat flux

The radiant heat flux of the Holuhraun lava flow field was estimated using MODIS (Moderate Resolution Imaging Spectroradiometer), elaborated by the MIROVA (Middle InfraRed Observation of Volcanic Activity) system of Coppola et al. (2016). During the effusive crisis, the MIROVA system allowed automatic measurements of the volcanic radiative power (VRP) sourced from the active lava flow field, and related effusion rates (Coppola et al., 2017). However, the VRP provided automatically by MIROVA is based on the simple heat flux conversion approach of Wooster et al. (2003), that works for active lava flows with high surface temperatures (>600 K). This approach begins to fail for cooling surfaces, especially below 200°C, which is the case of the post-eruption Holuhraun lava flow field. For this post-eruptive period we thus recalculate the radiant heat flux as:

\[
Q_{rad} = \sigma \varepsilon A_{flow} (BT_{MIR,flow}^4 - BT_{MIR,bk}^4)
\]

where \(\sigma\) is the Stefan-Boltzmann constant \((5.67051 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4})\), \(\varepsilon\) is the emissivity of the lava surface (assumed to be 0.95 for basalt; Patrick et al., 2004), \(A_{flow}\) is the final area of the cooling lava flow field \((84 \text{ km}^2;\) Pedersen et al., 2017), and \(BT_{MIR,flow}\) and \(BT_{MIR,bk}\) are the average brightness temperatures of the flow surface pixels and surrounding background respectively, calculated from selected MODIS-MIROVA cloud-free images (Figure 6).
Uncertainty and error

Individual SO₂ flux measurements were calculated as the product of the average vertical column density (VCD), effective haze width and wind speed. Thus, the total uncertainty on the flux measurement can be obtained by determining the uncertainties of these three independently derived quantities. The uncertainty on the average VCD depends on two factors: the intrinsic uncertainty of the retrieved SCDs at various angles and the uncertainty related to the method used to derive the mean VCD from a set of SCDs in one scan. The typical uncertainty of a DOAS SCD has been characterized for NOVAC instruments by Galle et al. (2010) and is expected to be around 15%. The sources of this uncertainty include added UV radiation by scattering between the plume and the instrument (Mori et al., 2006), uncertainty in the reference absorption cross sections used in the retrieval (Stutz and Platt, 1996), and possible distortions of the sensor response function due to temperature or other environmental effects (Galle et al., 2010).

The method proposed to derive a representative VCD is based on the assumption that a homogenous layer of gas surrounds the instrument. For each scan that contains a set of 26 SCDs taken at steps of 7.2° from horizon to horizon, this assumption results in a trough shape of the distribution of these SCDs. This occurs because the shortest column corresponds to the shortest path through the haze layer, when the scanner looks upwards, and then progressively increases with scan angle, as path length through the haze layer increases (see Figure 4). Thus, the measurements used in this study were selected by visual inspection of each scan, i.e. only scans with the expected trough-shaped distribution of SCDs were used for analysis. In addition, a numerical filter was applied to use only symmetrical scans, where the columns at +60° and -60° differed by less than 20%. A conservative uncertainty related to this method may be 20%,
allowing for imperfect scanning geometry, radiative transfer and natural variability of the
columns in the layer. Each VCD thus has a total uncertainty of ~35%.

The uncertainty related to the effective haze width depends on the stability of the wind direction
and accuracy in the determination of the dimensions of the lava flow field. The wind direction
used here was measured by a meteorological station in close proximity to the scanning system,
and these data were in good agreement with measurements from another station operated by
IMO 20 km further away. This suggests that wind directions close to ground level were
relatively stable over a large area, in spite of the presence of a hot lava flow field and a cold
glacier. The accuracy of the effective haze width for different wind directions must also be
considered. From a map of the extent of the lava flow field, the maximum widths of the lava
flow field in directions perpendicular to wind directions in the range 250° to 290° were
measured (Figure 5). This range of wind directions corresponds to the possible directions that
would produce a haze above the instrument, which was oriented at 270°. Because our study
took place after the eruption had ended, the extent of the lava flow field itself can be considered
unchanging. Assuming that the measurements of haze widths from a map can be achieved with
an accuracy of 1%, and that wind directions vary (spatially and temporarily) within 20%, the
total uncertainty on haze width is of the order of 20%.

Finally, uncertainty related to the haze speed must be determined. It is assumed here that haze
speed is equal to wind speed close to ground. We attribute a value of 20% to this uncertainty,
following Galle et al., (2002) and Edmonds et al., (2003). Therefore, the total uncertainty of a
single flux measurement is estimated to be, quite conservatively, 45%. The total uncertainty is
calculated by summing in “quadrature” the uncertainty of all variables, that is, the relative
uncertainty of the flux is equal to the square root of the sum of the squares of the relative
uncertainties of the variables, i.e. $\sqrt{(0.35^2+0.2^2+0.2^2)} = 0.45$. 

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Results

Lava flow field SO₂ flux

Post-eruption SO₂ fluxes from the Holuhraun lava flow field were calculated as varying between <1 kg/s and 9 kg/s (Figure 7). In the three months following the eruption, the lava flow field released an average of 3 kg/s of SO₂ (standard deviation: 1.9 kg/s). We found no overall trend in SO₂ emission with time, instead the SO₂ outgassing rate fluctuated without discernible trend during the three month period following the end of the eruption. The flux varied between 2 and 6 kg/s in the first weeks of March, then decreased and varied between 1 and 3 kg/s at the end of March. It became higher (~6 kg/s) and stable in April, then highly fluctuating again (1-9 kg/s) by the end of May 2015. The lack of DOAS scans with the required characteristics (trough-shaped distribution with wind direction blowing towards the instrument) in the first weeks of May and the scarcity of scans in April may reflect the less favourable weather conditions for useful DOAS measurements during the transition between winter and spring seasons.

DOAS data from the Holuhraun eruption were processed until the end of May 2015 (Figure 7), after which SO₂ from the lava flow field was below the detection limit of the instrument (5 ppm*m for single spectra). Thus, the lava flow field formed during a six-month long eruption continued to release measurable SO₂ for three months after the eruption ended.

Contribution of Holuhraun lava flow field SO₂ emissions
Assuming that the average post-eruption flux of 3 kg/s was emitted constantly throughout the three months of this study, an estimated total of 24 kt of SO$_2$ was released during this period. This is equivalent to an additional 0.2% to the total of ~11 Mt of SO$_2$ that was released during the eruption. This estimated 24 kt of SO$_2$, released during three months, is greater than the 16 kt of SO$_2$ emitted by Icelandic industry in 2013 (Centre on Emission Inventories and Projections, 2015). The SO$_2$ emitted by the lava flow field remained near ground level, resulting in high ground-level concentrations of SO$_2$ near to the eruption site, and posing potential health hazards (including respiratory problems and eye irritation; Longo et al, 2008; Ilyinskaya et al., 2017) to people exposed to it. Access was therefore restricted to the area around the lava flow field until 1 June 2015, after a field visit by IMO on 19 May 2015 measured only minor emissions of SO$_2$ from fractures in the lava and from the main eruption crater (IMO, 2015c).

The Holuhraun eruption magma was highly rich in sulphur (Gauthier et al., 2016; Gíslason et al., 2015), resulting in high emissions of SO$_2$ at the vent. It is likely that the magma had a permeable bubble network during ascent allowing volatiles to be released from the magma to form the gas-rich eruption plume (e.g. Burton et al., 2007; Polacci et al., 2008). Therefore, when the lava was erupted much of its sulphur content had already been released by efficient degassing at the vent, and so only a small proportion of SO$_2$ remained within the lava to be released during and after emplacement. Due to the large size of this flood basalt eruption, this small proportion of SO$_2$ released by the lava flow field was still a large and significant mass of SO$_2$.

The petrologic approach of Bali et al. (submitted) was used to calculate the portion of erupted lava responsible for outgassing the post-eruptive SO$_2$. Degassed tephra glass contained 377 ppmw sulphur (Bali et al.), and degassed lava contained 97 ppmw sulphur (Gauthier et al.,
The 24 kt of post-eruptive SO$_2$ are calculated to have been emitted from $4.23 \times 10^{10}$ kg of the $3.56 \times 10^{12}$ kg of lava produced by the Holuhraun eruption.

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**Lava flow field radiant heat flux**

The radiant heat flux due to the Holuhraun lava flow field decreased following the end of the eruption (Figure 8), from $13.3 \times 10^3$ MW on 8 February 2015 (in the final month of the eruption) to $6.9 \times 10^3$ MW on 10 March (10 days after the eruption ended), followed by a continued cooling. Once the eruption had ended, the source of hot lava was removed, and so all heat released was due to the cooling of the existing lava. The measured post-eruption SO$_2$ fluxes were emitted as this emplaced lava cooled.

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**Discussion**

**Cooling and outgassing of lava flows**

Several studies have analysed the cooling and crystallisation of lava bodies (e.g. Keszthelyi and Denlinger, 1996; Wooster et al., 1997; Cashman et al., 1999; Patrick et al., 2004; Kolzenburg et al., 2016). Lava flows cool by thermal radiation and atmospheric convection from the surface, and by conduction to the underlying country-rock (Keszthelyi and Denlinger, 1996; Neri, 1998; Patrick et al., 2004). Lava from the Holuhraun eruption had up to 45% vesicles (Lavallée et al., 2015). This high vesicularity is likely to have enhanced the cooling of the flows as lava porosity is considered to affect the rate of cooling, with more vesicular lava cooling more rapidly (Keszthelyi and Denlinger, 1996).
Processes that occur during cooling could contribute to the release of SO$_2$ from the lava flow. As an emplaced lava flow cools and solidifies post-eruption, it thermally contracts (Spry, 1962; Wittmann et al., 2017). Thermal contraction of lava results in the formation of cooling fractures such as brittle cracks and columnar jointing (Ryan and Sammis, 1978; Long and Wood, 1986; Reiter et al., 1987; Grossenbacher and McDuffie, 1995). Continued cooling causes fractures to propagate from the surface into the interior of the flow (Reiter et al., 1987; Aydin and DeGraff, 1988; Kattenhorn and Schaefer, 2008). Not only will fractures and cracks enhance the cooling of the lava flow interior, they will also provide pathways for volatiles to be released from the interior of the flow to the atmosphere (Fuller, 1938; Wilmoth and Walker, 1993).

Vesicles within a lava flow consist of both a primary population that were exsolved from the melt during eruption, and a secondary population that were exsolved during cooling and crystallisation (Cashman et al., 1994). The strength of the cooling lava crust will prevent the upward migration of bubbles, causing volatiles to remain trapped within the flow (Polacci and Papale, 1997). Volatiles exsolved from the lava therefore contribute to the formation of a low permeability vesicle network in which the volatiles become trapped, with inefficient outgassing due to the formation of a solid outer crust on the lava flow. The outgassing efficiency will be dependent on the permeability of the lava flow (Burgisser and Degruyter, 2015; Kennedy et al., 2016), with the permeability dependent on vesiculation history (Saar and Manga, 1999; Wright et al., 2009). Fracturing of the lava will then create high permeability pathways allowing efficient outgassing of these trapped volatiles, as seen at the Holuhraun lava flow field (Figure 9).

Extraction of gas from the lava flow should be most efficient where it is highly fractured, which will be where the thermal contraction is greatest. The formation of high permeability pathways through propagation of cooling fractures would account for the fact that SO$_2$ was released from
the Holuhraun lava flow field for several months after lava flow emplacement, and that over this time the SO$_2$ flux did not decrease concurrently with the decrease in thermal emission. SO$_2$ flux is not a simple decay correlated with lava cooling, as both volatile availability and fracture propagation produce the SO$_2$ flux. As fractures provide a pathway for gas to escape, outgassing is able to continue as fracture propagation continues to tap volatiles trapped in the cooling flow interior. The episodic nature of fracture propagation (Ryan and Sammis, 1978; DeGraff and Aydin, 1987; Reiter et al., 1987) would explain the fluctuating post-eruption SO$_2$ flux.

Using a petrological method, Thordarson et al. (1996) calculated that of the SO$_2$ released by the 1783-1784 Laki eruption, 20% was released by lava flows during their emplacement and cooling. This is two orders of magnitude larger than our calculated value of 0.2% of SO$_2$ released by the Holuhraun lava flow field in the three months following the eruption. The reason for this large discrepancy is that our measurements only began after the eruption had ended. The Holuhraun eruption lasted for six months, and therefore lava flows had been outgassing for up to six months before our study even began, and so a significant proportion of the outgassing has not been captured. Outgassing from lava flows is likely to have been greatest during emplacement, when mechanical break up of the lava crust facilitated gas release (Polacci and Papale, 1997; Soule and Cashman, 2004), and we did not capture syn-emplacement outgassing during our study.

The hazard caused by the outgassing of the Holuhraun lava flow field remained at ground level as SO$_2$ formed a low-lying haze, and it continued to occur after the eruption ended as the lava flow field continued to outgas.

Conclusions
Remote sensing scanning DOAS instruments were installed at the eruption site of the 2014-2015 Holuhraun eruption to monitor the flux of SO$_2$. Based on a novel approach to measure emissions from the emplaced lava flow field, we found that this field continued to release SO$_2$ for three months after the end of the eruption. During this time, SO$_2$ flux from the lava flow field varied from <1 kg/s to 9 kg/s, and a post-eruption average flux of 3 kg/s was calculated. The post-eruption Holuhraun lava flow field was a significant source of environmental pollution, releasing 24 kt of SO$_2$ in the three months after the eruption. This post-eruption lava flow field SO$_2$ flux contributed an additional 0.2% to the ~11 Mt total eruption SO$_2$ emissions in the three months following the eruption. This emission remained near ground level, posing health hazards, which resulted in access to the area being restricted until 1 June 2015, three months after the eruption ended. During solidification, SO$_2$ can be outgassed, we propose, by fracture propagation as high permeability pathways are developed, allowing gas to be released for several months after the end of the eruption. We show that this release is not a simple decay correlated with lava cooling, as both volatile availability and fracture propagation produce a fluctuating SO$_2$ flux.

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Figures and captions

Fig. 1  a) Map of Iceland showing the location of the eruption (triangle).  b) The location of the 2014-2015 Holuhraun eruption site. The fissure is represented by the dashed line and the area of the lava flow field is highlighted. The location of ScanDOAS instrument MAYP111126 is indicated by the symbol x. The viewing direction (arrow) and scanning sector (arc segment) of the instrument are shown. Base maps from the Icelandic Meteorological Office (IMO), 2015b
Fig. 2  Two-stage degassing during basaltic fissure eruptions: $\text{SO}_2$ is released at vents contributing to an eruption plume and from lava flows forming a low-level haze. Adapted from Thordarson and Self, 2003. Inset images show a) the Holuhraun eruption plume and b) gas released by the lava flow field. Photos by B. Bergsson (IMO)
Fig. 3 A photo taken during the surveillance flight on 4th November 2014 showing the low-level gas haze and the higher-level eruption plume. Gases (including H$_2$O, both magmatic and meteoric, SO$_2$ and CO$_2$) can be seen rising from fractures in the lava flow field. Photo by M. Hensch, IMO
Fig. 4 A schematic illustration of the characteristic shapes recorded by scanning DOAS when scanning vertically through a) an elevated eruption plume produced by SO$_2$ released at a vent, and b) a low-level haze produced by SO$_2$ released from a lava flow. c) DOAS scan from 11:57 on 7th March 2015 showing the characteristic trough shape of a scan made through the low-level SO$_2$ haze released by the Holuhraun lava flow field. Black columns are the SO$_2$ column densities at incremental angles.
Fig. 5 The effective plume width ($L$) is the maximum width of the lava flow field in the direction perpendicular to the wind direction. The diagram shows the effective plume widths (dashed lines) for wind directions (arrows) of 250°, 270° (directly towards the viewing direction of the ScanDOAS instrument) and 290°. The location of ScanDOAS instrument MAYP111126 is indicated by the symbol x.
**Fig. 6** A MODIS-MIROVA thermal map (1 km resolution) elaborated for radiant flux calculation during the cooling phase. The map represents the brightness temperature (in K) recorded by the MIR channel of MODIS (centred at ~3.9 μm). Note the temperature anomaly due to the presence of the Holuhraun lava flow field to the south of Askja caldera.
Fig. 7 Post-eruption SO$_2$ flux from the 2014-2015 Holuhraun lava flow field as recorded by ScanDOAS instrument MAYP111126. Labelled dates are at 10 day intervals, error bars are 45%
Fig. 8 Radiant heat flux calculated during the post-eruptive cooling phase of the Holuhraun lava flow field according to equation 4. The vertical dashed line indicates the end of the eruption on 27th February 2015. Labelled dates are at 50 day intervals.
Fig. 9  A photo taken during a surveillance flight on 4th November 2014 showing gases being released from cooling fractures in the lava flow field (highlighted by arrows). Gases include H$_2$O (both magmatic and meteoric), SO$_2$ and CO$_2$. Photo by M. Hensch, IMO.