Improving the security of quantum direct communication with authentication

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Two protocols of quantum direct communication with authentication [Phys. Rev. A 73, 042305 (2006)] are recently proposed by Lee, Lim and Yang. In this paper we will show that in the two protocols the authenticator Trent should be prevented from knowing the secret message of communication. The first protocol can be eavesdropped by Trent using the intercept-measure-resend attack, while the second protocol can be eavesdropped by Trent using single-qubit measurement. To fix these leaks, I revise the original versions of the protocols by using the Pauli-Z operation $\sigma_z$ instead of the original bit-flip operation $X$. As a consequence, the protocol securities are improved.

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Quantum key distribution (QKD) is one of the most interesting topics in quantum information processing, which provides a novel way for two legitimate parties to share a common secret key over a long distance with negligible leakage of information to an eavesdropper Eve. Its ultimate advantage is the unconditional security. Hence, after Bennett and Brassard’s pioneering work published in 1984[1], much attentions have been focused on this topic and a variety of quantum communication protocols[1-13,17-18] have been proposed. In these works, various properties of quantum mechanics, such as no-cloning theorem, uncertainty principle, entanglement, indistinguishability of non-orthogonal states, non-locality, and so on, are used to accomplish QKD tasks. Different from QKD, the deterministic quantum secure direct communication (QSDC) protocol is to transit directly the secret messages without first generating QKD to encrypt them. Hence it is very useful and usually desired, especially in some urgent time. However, a deterministic secure direct communication protocol is more demanding on the security. Therefore, only recently a few of deterministic secure direct communication protocols have been proposed[3-12] and some of them are essentially insecure[13-15]. Two of the QSDC protocols are the Lee-Lim-Yang protocols of quantum direct communication with authentication[12] proposed very recently by Lee, Lim and Yang using the Greenberg-Horne-Zeilinger (GHZ) states[16]. Based on some security analysis Lee, Lim and Yang claimed that their two protocols are secure. However, in this paper we will show that in the two protocols the authenticator Trent should be prevented from knowing the secret message of communication. The first protocol can be eavesdropped by Trent using the intercept-measure-resend attack, while the second protocol can be eavesdropped by Trent using single-qubit measurement. I will fix these leaks by modifying the original versions of the protocols using that Pauli-Z operation $\sigma_z$ instead of the original bit-flip operation $X$. As a consequence, the protocol securities are improved.

There are three parties in each of the Lee-Lim-Yang protocols. Alice and Bob are the legitimate users of the communication. Trent is the third party who is introduced to authenticate the two users participating in the communication. Trend is assumed to be more powerful than the other two parties and he supplies the GHZ states each in the form of $|\Psi\rangle = (|000\rangle + |111\rangle)/\sqrt{2}$. The protocols are composed of two parts: one is for an authentication process and the other is for a direct communication. The authentication process is same for both Lee-Lim-Yang protocols. After the authentication, there are two possibilities for Alice to send qubits: one is to Bob and the other is to Trend. The former
case corresponds to the Lee-Lim-Yang protocol 1 and the latter case to the Lee-Lim-Yang protocol 2. This is a difference between the two protocols. Incidentally, there is an unphysical mistake about the correspondences in the text of Ref.[12].

The purpose of the authentication process in the Lee-Lim-Yang protocols is to let the three participants safely share GHZ states. To achieve this goal, it is assumed that Trend should share in priori secret authentication keys $K_{ta}$ and $K_{tb}$ with Alice and Bob, respectively. The lengths of the authentication keys $K_{ta}$ and $K_{tb}$ are larger than the length of the bit string of the secret message which will be communicated from Alice to Bob. According to the one-time pad cryptography, when the private key length is equal to the secret message length, the secret message can be securely communicated to remote places after encryption. If Trend is not assumed to prevent from knowing the secret message, then in this case the secret message can be transferred in the following very simple classical way instead of using the Lee-Lim-Yang QSDC protocols, i.e, Alice can securely send the secret message to Trend by using their secret authentication key $K_{ta}$ and then Trend can securely send the secret message to Bob by using their secret authentication key $K_{tb}$. Hence, in the Lee-Lim-Yang protocols the third party Trend should be prevented from knowing the secret message though he is introduced to authenticate the communication.

Assume that the GHZ states are safely shared among the three parties after the authentication process. Let us now briefly review the second part of the Lee-Lim-Yang protocol 1.

(a) Alice selects a subset of GHZ states of her remaining set after authentication and keeps it secret.
(b) Alice chooses a random bit string which has no correlation to the secret message to transmit to Bob. This bit string will be used to check the security of the channel.
(c) Following the random bit string, Alice performs unitary operations on the the qubits selected for check process. The unitary operations are defined as follows. The bit ’0’ correspond to $H$ and the bit ‘1’ to $HX$, where $H$ is the Hadamard operation and $X$ is the bit-flip operation. The GHZ states after Alice’s operations are transformed into:

$$H_{A}|\Psi\rangle = H_{A}(|000\rangle_{ATB} + |111\rangle_{ATB})/\sqrt{2}$$

$$= \frac{1}{2}(|000\rangle_{ATB} + |100\rangle_{ATB} + |011\rangle_{ATB} + |111\rangle_{ATB})$$

$$= \frac{1}{2}\{|\phi^{+}\rangle_{AB} - |\psi^{-}\rangle_{AB}\rangle - T + (|\phi^{-}\rangle_{AB} + |\psi^{+}\rangle_{AB}\rangle + T\}$$

(1)

$$H_{AX_{A}}|\Psi\rangle = H_{A}(|100\rangle_{ATB} + |011\rangle_{ATB})/\sqrt{2}$$

$$= \frac{1}{2}(|000\rangle_{ATB} - |100\rangle_{ATB} + |011\rangle_{ATB} + |111\rangle_{ATB})$$

$$= \frac{1}{2}\{|\phi^{-}\rangle_{AB} - |\psi^{+}\rangle_{AB}\rangle - T + (|\phi^{+}\rangle_{AB} + |\psi^{-}\rangle_{AB}\rangle + T\}$$

(2)

where $|\phi^{\pm}\rangle = (|00\rangle \pm |11\rangle)/\sqrt{2}$, $|\psi^{\pm}\rangle = (|01\rangle \pm |10\rangle)/\sqrt{2}$ and $|\pm\rangle = (|0\rangle \pm |1\rangle)/\sqrt{2}$.

d) Alice encodes the secret message with a classical error correction code on the remaining GHZ states in terms of the unitary operation definition in (c).

(e) After making all unitary operations, Alice sends the encoded qubits to Bob.

(f) Bob makes Bell measurements on the pairs of particles consisting of his qubits and Alice’s qubits.

(g) Trent measures his third qubit in the $x$ basis $\{|-, +\rangle\}$ and publishes the measurement outcomes.

(h) Using Trend’s measurement outcomes and his Bell-state measurement outcomes, Bob infers Alice’s secret bits consisting of both the random bits and the secret message.

(i) Bob lets Alice reveal the check bits’ positions and values.

(j) Bob can know whether the channel is disturbed according to the error rate. If the error rate is higher than expected, an eavesdropper is concluded in the communication but fortunately the secret
message is not leaked out. If the error rate is lower, Bob can extract the secret message (see Table I in Ref.[12]).

As I show before, the authenticator Trent should be prevented from the secret message. Otherwise, the communication can be realized in a very simple classical way. Although the Lee-Lim-Yang protocol 1 is claimed to be secure (that is, the secret message cannot be leaked out), in the protocol 1 the insider Trent can eavesdrop the secret message by using the intercept-measure-resend attack. This can be seen as follows. When Alice sends her encoded qubits to Bob, Trent intercepts the qubits and performs $H$ operation on each qubit. In this case, the whole states are transformed into

$$H_A H_A |\Psi\rangle = (|000\rangle_{ATB} + |111\rangle_{ATB})/\sqrt{2}. \tag{3}$$
$$H_A H_A X_A |\Psi\rangle = (|100\rangle_{ATB} + |011\rangle_{ATB})/\sqrt{2}. \tag{4}$$

After the unitary operations, Trent measures Alice’s qubit and his qubit in the $z$ basis $\{|0\rangle, |1\rangle\}$, respectively. If the two outcomes are same, then Trent can conclude that Alice has performed a $H$ operation corresponding to the bit ‘0’. Otherwise, Alice has performed a $HX$ operation corresponding to the bit ‘1’. In this case, Trent has already known for each Alice’s qubit which unitary operation Alice has performed on. Alternatively, he has already got Alice’s whole bit string including both the random bit string for check and the secret message. In the following what he needs to do is to remove the random bits. Fortunately, in the step (i), Alice will publish which qubits are used as check qubits. This means that the authenticator Trent can completely know the secret message using this intercept-measure-resend attack.

After Trent’s attack, he sends Alice’s qubits to Bob. One can easily find that for Alice and Bob the error rate will obviously be higher than expected. Unfortunately, Alice and Bob only know the channel is disturbed and still believe that the secret message is not leaked out.

To fix this leak, we think the original version of Lee-Lim-Yang protocol 1 can be modified by using the Pauli-Z operation $\sigma_z$ instead of the original bit-flip operation $X$. In this case, the total states after Alice’s operations are represented as follows:

$$H_A |\Psi\rangle = H_A(|000\rangle_{ATB} + |111\rangle_{ATB})/\sqrt{2}$$
$$= \frac{1}{2}(|000\rangle_{ATB} + |100\rangle_{ATB} + |011\rangle_{ATB} - |111\rangle_{ATB})$$
$$= \frac{1}{2} \{ (|\phi^+\rangle_{AB} - |\psi^-\rangle_{AB})|0\rangle_T + (|\phi^-\rangle_{AB} + |\psi^+\rangle_{AB})|1\rangle_T \}, \tag{5}$$

$$H_A \sigma_z A |\Psi\rangle = H_A(|000\rangle_{ATB} - |111\rangle_{ATB})/\sqrt{2}$$
$$= \frac{1}{2}(|000\rangle_{ATB} + |100\rangle_{ATB} - |011\rangle_{ATB} + |111\rangle_{ATB})$$
$$= \frac{1}{2} \{ (|\phi^-\rangle_{AB} + |\psi^+\rangle_{AB})|0\rangle_T + (|\phi^+\rangle_{AB} - |\psi^-\rangle_{AB})|1\rangle_T \}. \tag{6}$$

After this modification, if Trent intercepts Alice’s encoded qubits and performs $H$ operations, then the total states are transformed into

$$H_A H_A |\Psi\rangle = (|000\rangle_{ATB} + |111\rangle_{ATB})/\sqrt{2}, \tag{7}$$
$$H_A H_A \sigma_z A |\Psi\rangle = (|000\rangle_{ATB} - |111\rangle_{ATB})/\sqrt{2}. \tag{8}$$

If he measures respectively his qubit and Alice’s qubit in the $z$ basis, the outcomes will always be same. In this case, he can not know for each of Alice’s qubits which unitary operation Alice has performed on. This means Trent can not know Alice’s secret message. However, one can easily find that revised protocol works successfully as far as Alice and Bob’s communication is concerned. See Table 1 for a brief summary.
Table 1 The relation of Alice’s operation, Bob’s measuremnt, and Trent’s announcement in the revised Lee-Lim-Yang protocol 1 can be summarized as follows.

| Trent’s publication | Bob’s Measurement | Alice’s operation |
|---------------------|-------------------|------------------|
| $|+\rangle_T$       | $|\phi^+\rangle_{AB}$ or $|\psi^+\rangle_{AB}$ | $H\sigma_z$ (1) |
| $|+\rangle_T$       | $|\phi^-\rangle_{AB}$ or $|\psi^-\rangle_{AB}$ | $H$ (0)          |
| $|-\rangle_T$       | $|\phi^+\rangle_{AB}$ or $|\psi^-\rangle_{AB}$ | $H$ (0)          |
| $|-\rangle_T$       | $|\phi^-\rangle_{AB}$ or $|\psi^+\rangle_{AB}$ | $H\sigma_z$ (1) |

Let us now briefly review the second part of the Lee-Lim-Yang protocol 2. Some steps are same for both protocols. Nevertheless, for completeness, I depict all the steps as follows.

(a’) Alice selects a subset of GHZ states of her remaining set after authentication and keeps it secret.

(b’) Alice chooses a random bit string which has no correlation to the secret message to transmit to Bob. This bit string will be used to check the security of the channel.

(c’) Following the random bit string, Alice performs unitary operations on the the qubits selected for check process. The GHZ states after Alice’s operations are transformed into:

$$H_A|\Psi\rangle = H_A(|000\rangle_{AB} + |111\rangle_{AB})/\sqrt{2}$$

$$= \frac{1}{2} \{(|\phi^+\rangle_{AT} - |\psi^-\rangle_{AT})|+\rangle_B + (|\phi^-\rangle_{AT} + |\psi^+\rangle_{AT})|+\rangle_B\}.$$  

$$H_A X_A |\Psi\rangle = H_A(|100\rangle_{AB} + |011\rangle_{AB})/\sqrt{2}$$

$$= \frac{1}{2} \{(|\phi^-\rangle_{AT} - |\psi^+\rangle_{AT})|+\rangle_B + (|\phi^+\rangle_{AT} + |\psi^-\rangle_{AT})|+\rangle_B\}. \quad (10)$$

(d’) Alice encodes the secret message with a classical error correction code on the remaining GHZ states in terms of the unitary operation definition in (c’).

(e’) After making all unitary operations, Alice sends the encoded qubits to Trent.

(f’) Trent makes Bell measurements on the pairs of particles consisting of his qubits and Alice’s qubits and publishes the measurement outcomes.

(g’) Bob measures his qubits in the $x$ basis $\{|-\rangle, |+\rangle\}$.

(h’) Using Trend’s Bell-state measurement outcomes and his measurement outcomes, Bob infers Alice’s secret bits consisting of both the random bits and the secret message.

(i’) Bob lets Alice reveal the check bits’ positions and values.

(j’) Bob can know whether the channel is disturbed according to the error rate. If the error rate is higher than expected, an eavesdropper is concluded in the communication but fortunately the secret message is not leaked out. If the error rate is lower, Bob can extract the secret message (see Table II in Ref.[12]).

In the protocol 2 the insider Trent can eavesdrop the secret message by using the measurement attack as follows. Trent performs $H$ operation on each qubit he received from Alice. In this case, the whole states are transformed into

$$H_A H_A |\Psi\rangle = (|000\rangle_{AB} + |111\rangle_{AB})/\sqrt{2}, \quad (11)$$

$$H_A H_A X_A |\Psi\rangle = (|100\rangle_{AB} + |011\rangle_{AB})/\sqrt{2}. \quad (12)$$

After the unitary operations, Trent measures Alice’s qubit and his qubit in the $z$ basis $\{|0\rangle, |1\rangle\}$, respectively. If the two outcomes are same, then Trent can conclude that Alice has performed a $H$ operation corresponding to the bit ‘0’. Otherwise, Alice has performed a $HX$ operation corresponding to the bit ‘1’. In this case, Trent has already known for each Alice’s qubit which unitary operation Alice has performed on. Alternatively, he has already got Alice’s whole bit string including both the random bit
string for check and the secret message. In the following what he needs to do is to remove the random bits. Fortunately, in the step (i'), Alice will publish which qubits are used as check qubits. This means that the authenticator Trent can completely know the secret message using this measurement attack.

After Trent’s attack, he randomly publishes his measurement outcomes. One can easily find that for Alice and Bob the error rate will obviously be higher than expected. Unfortunately, Alice and Bob only know the channel is disturbed and still believe that the secret message is not leaked out.

To fix this leak, the original version of Lee-Lim-Yang protocol 2 can also be modified by using the Pauli-Z operation $\sigma_z$ instead of the original bit-flip operation $X$. In this case, the total states after Alice’s operations are represented as follows:

$$H_A|\Psi\rangle = H_A((000)_{AB} + (111)_{AB})/\sqrt{2}$$
$$= \frac{1}{2}\{(|\phi^+\rangle_{AB} - |\psi^-\rangle_{AB})|\rangle_T + (|\phi^-\rangle_{AB} + |\psi^+\rangle_{AB})|+\rangle_T\}. \quad (13)$$

$$H_A\sigma_z|\Psi\rangle = H_A((000)_{AB} - (111)_{AB})/\sqrt{2}$$
$$= \frac{1}{2}\{(|\phi^-\rangle_{AB} + |\psi^+\rangle_{AB})|\rangle_B + (|\phi^+\rangle_{AB} - |\psi^-\rangle_{AB})|+\rangle_B\}. \quad (14)$$

After this modification, if Trent intercepts Alice’s encoded qubits and performs $H$ operations, then the total states are transformed into

$$H_AH_A|\Psi\rangle = (000)_{AB} + (111)_{AB})/\sqrt{2}, \quad (15)$$

$$H_AH_A\sigma_z|\Psi\rangle = (000)_{AB} - (111)_{AB})/\sqrt{2}. \quad (16)$$

If he measures respectively his qubit and Alice’s qubit in the $z$ basis, the outcomes will always be same. In this case, he can not know for each of Alice’s qubits which unitary operation Alice has performed on. This means Trent can not know Alice’s secret message. However, one can easily find that revised protocol works successfully as for as Alice and Bob’s communication is concerned. See Table 2 for a brief summary.

**Table 2** The relation of Alice’s operation, Bob’s measurement, and Trent’s announcement in the revised Lee-Lim-Yang protocol 2 can be summarized as follows.

| Trent’s announcement | Bob’s Measurement | Alice’s operation |
|----------------------|------------------|------------------|
| 0 $(|\phi^+\rangle_{AB}$ or $|\psi^-\rangle_{AB})$ | $|+\rangle_T$ | $H\sigma_z$ (1) |
| 0 $(|\phi^+\rangle_{AB}$ or $|\psi^-\rangle_{AB})$ | $|-\rangle_T$ | $H$ (0) |
| 1 $(|\phi^-\rangle_{AB}$ or $|\psi^+\rangle_{AB})$ | $|+\rangle_T$ | $H$ (0) |
| 1 $(|\phi^-\rangle_{AB}$ or $|\psi^+\rangle_{AB})$ | $|-\rangle_T$ | $H\sigma_z$ (1) |

To summarize, in this paper we have shown that the Lee-Lim-Yang protocols can be eavesdropped by the authenticator Trent using some specific attacks and we have revised the original versions of the protocols by using the Pauli-Z operation $\sigma_z$ instead of the original bit-flip operation $X$ so that the protocol securities are improved.

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