Vibrational Relaxation of Functional Groups in dAMP Molecules Probed with UV Resonance Raman Spectroscopy

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Dynamical properties of functional groups in 2’-deoxyadenosine-5’-monophosphate (dAMP) compound, were identified by UV resonance Raman spectroscopy (UVRR), upon varying nucleotide concentration in aqueous solution (200-600 μM). The studied full-widths at half-maximum (fwhm’s) were found between 13 - 21 cm⁻¹ and the corresponding global relaxation times were faster than 0.817 ps and slower than 0.506 ps. Also, the band around 1430 cm⁻¹ (C4N9-O5C8H) in the UV resonance Raman spectrum of dAMP molecule at 400 μM concentration in aqueous solution, was selected for vibrational band shape analysis through time correlation function (CF) concept. Current theories developed for vibrational dephasing (Kubo-Rothschild and Oxtoby) have been applied to this profile and relevant relaxation parameters have been obtained and discussed. The best fit parameters for this dissipation channel of the vibrational excitation energy were established. To our knowledge this is the first UVRR study on nucleotide vibrational band shape analysis through time correlation function concept.

Keywords: 2’-deoxyadenosine-5’-monophosphate, UV resonance Raman spectroscopy, time correlation function, vibrational relaxation, vibrational dephasing

Nucleic acids components play an important role in different stages of the biological processes. Inside the living cell, the static structures of nucleic acids and their constituents must function in dynamic processes, through participation in specific interactions with other molecules. The dynamics of nucleic acids and their components is crucial for biological function of these biomolecules, particularly playing a significant role in the molecular recognition of DNA-protein and DNA-ligand systems [1]. It is clear that the understanding of the complex processes of DNA biopolymer in vivo requires the detailed knowledge of its dynamics as a function of the molecular relaxation characterizing its components.

Insight knowledge about the molecular dynamics processes involved in liquids mixtures can be obtained by Raman spectroscopy [2]. However, a simple analysis of the vibrational half bandwidths is insufficient for comprehending the dynamics of a molecule in liquid. Therefore, in the case of Raman profiles, it is necessary to pass from the spectral into a time range and to consider the time correlation functions (CFs) of vibrational and orientational relaxation, the Fourier transforms of band contours detected in the Raman spectra. From this stand point, Rabadanov et al. [3] studied the dephasing of totally symmetrical vibration ν₁(A₁) and overtone ν₂(E) of the perchlorate ion in molten lithium perchlorate. Their findings show that vibrational dephasing is the main factor of the broadening of isotropic spectral bands and in this particular case, dephasing behaves according to the Gaussian-Markovian mechanism [3].

The sensitivity of vibrational spectroscopy to intermolecular forces and its potential to probe solute-solvent interactions were investigated by Ramakrishnan et al. in substituted benaldehyde binary mixtures [4]. In this context, the vibrational CF of the ν(C≡N) mode, proved the presence of a dynamical regime in the system, positively influenced by the temperature increase [5].

Most recently, studies concerning the role of microenvironment and hydrodynamic forces or of the dilution effect on the vibrational relaxation processes were carried out by Kalamponias et al. on pyridine-water complexes [6] and on PbCl₂-KCl ionic liquids [7], respectively. By taking into consideration vibrational second moments variation, dephasing times or the vibrational relaxation rate as guiding parameters, they concluded that the microenvironment plays an important role in the vibrational relaxation processes (in correlation to solvent viscosity, density, refractive index, etc.).

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In particular, advances in this research field are contributing to build an interesting picture of the molecular dynamics of nucleic acids, their derivatives and components, e.g. the study carried out by Kozich et al. [8]. The symmetric stretching mode, νs(PO$_2^-$), of several mononucleotides (5'-GMP, 5'-CMP and 5'-dCMP) have been studied in H$_2$O and D$_2$O solutions, at different concentrations and temperatures by infrared spectroscopy [9-12]. Studies on the fwhm’s of calf-thymus DNA Raman bands revealed their sensitivity to pH, as well as to the presence and concentration of different metal cations [13]. Besides, our group investigated the vibrational relaxation of a 22-mer duplex DNA, as probed with UVRR spectroscopy [14]. Resonance Raman spectroscopy (RRS) is a useful technique employed to obtain experimental data on molecular relaxation dynamics [15].

All these research findings underline the complex nature of the vibrational relaxation processes in different molecular systems.

In the followings, we report an UV resonance Raman spectroscopic study on the subpicosecond dynamics of functional groups in 2'-deoxyadenosine-5'-monophosphate (dAMP) in aqueous solutions, at different concentrations. We will concentrate on the full-widths at half-maximum (fwhm’s) and on the corresponding global relaxation times. The existence of molecular concentration-dependent processes with biological relevance, justifies the study of biomolecular dynamics upon variation of this parameter. Also, square root of concentration-dependences of the half bandwidths of different vibrational modes, can give information on the most efficient relaxation mechanism which takes place in aqueous system [11, 16-20]. Besides, the 1430 cm$^{-1}$ UVRR band characterizing ν(C4N9-δC8H) oscillator of dAMP molecule (400 μM), was selected for band shape analysis. Vibrational dephasing accomplishes an efficient vibrational relaxation of this mode. Theoretical models (Kubo-Rothschild and Oxtoby) for this process have yielded significant relaxation parameters.

**Experimental part**

2'-deoxyadenosine-5'-monophosphate (D-6375 Sigma) has been used without further purification. dAMP molecules were dissolved in aqueous solutions (10 mM Bistris, 50 mM NaCl, pH 6.4), at different concentrations: 200 μM, 300 μM, 400 μM, 500 μM and 600 μM (2·10$^{-4}$, 3·10$^{-4}$, 4·10$^{-4}$, 5·10$^{-4}$, 6·10$^{-4}$ mol· dm$^{-3}$).

The UV resonance Raman measurements were done at the University of Osnabrück, Osnabrück, Germany. Continuous wave Raman equipment described in detail elsewhere [14,21] was employed for recording the UVRR spectra. All UV resonance Raman profiles of dAMP molecules were registered in the same experimental conditions.

**Data treatment**

Data treatment was applied as described in Muntean et al. [14].

**Results and discussions**

In this work the subpicosecond dynamics of dAMP functional groups, in the case of relaxation processes which take place in solution are of interest.

Molecular structure of 2'-deoxyadenosine-5'-monophosphate is presented elsewhere [22]. The 275 nm-excited UV resonance Raman spectra of dissolved dAMP molecules, at different concentrations, are presented in Figure 1. Spectra-structure correlations were established for the UV resonance Raman profiles of the investigated compound [23,24].

**dAMP concentration-dependences of full-widths at half-maximum and global relaxation times**

The selection of the concentration in the interval 2·10$^{-4}$ - 6·10$^{-4}$ mol· dm$^{-3}$ is depending on the hypothesis to work with an isolated molecule, i.e. the relaxation processes are related to an isolated molecule present in the aqueous medium. Besides, appropriated experimental conditions have been selected in relation with the hypothesis established in the theoretical models, in order to apply the vibrational dephasing models (diluted concentrations to avoid intermolecular interactions). A possible aggregation process of the nucleotide at concentrations higher than ~0.3 mol dm$^{-3}$ was previously detected (refs [11-12] and references therein).
Fig. 1. The 275 nm-excited UV resonance Raman spectra of dissolved 2’-deoxyadenosine-5’-monophosphate molecules in 10 mM Bistris, 50 mM NaCl, pH 6.4, at different concentrations: A. 200 μM; B. 300 μM; C. 400 μM; D. 500 μM; E. 600 μM.

Assuming a Lorentzian profile for the experimental band shape, a condition very often realized in solution, the full-widths at half-maximum (fwhm) may be expressed by:

$$ \Delta \nu_{1/2} = (c \pi \tau)^{1/2} $$  \hspace{1cm} (1)

where $\tau$ is a “relaxation time” relative to all the relaxation processes contributing to the observed band profile, and $c$ is the velocity of light [12,25].

Dynamical parameters associated with the investigated subpicosecond dynamics are presented in Table 1. Vibrational UVRR full-widths at half-maximum studied for the functional groups of dAMP molecules are to be found in the
wavenumber range 13 - 21 cm\(^{-1}\) and the corresponding global relaxations times \(\tau\) are faster than 0.817 ps and slower than 0.506 ps. In the case of dAMP relaxation processes in aqueous solutions, taking into account the values of the time dynamic parameter (\(\tau_{UVRR}\)), we can suppose that the dominant relaxation mechanism is the vibrational one.

Analyzing data in Table 1 on dAMP concentration, several remarks are to be done. For the band near 1318 cm\(^{-1}\) (N9C8+N3C2+\(\delta\)C8H-\(\delta\)C2H), global relaxation time increases between 200-400 \(\mu\)M and decreases in the interval 500-600 \(\mu\)M. In the case of 1433 cm\(^{-1}\) (C4N9-\(\delta\)C8H) vibrational profile, global relaxation time is equal to 0.758 ps at 200 \(\mu\)M, 400 \(\mu\)M and 600 \(\mu\)M concentrations. As far as the mode near 1494 cm\(^{-1}\) (N9C8+\(\delta\)C8H) is concerned, the global relaxation time is equal to 0.685 ps at 200 \(\mu\)M, 400 \(\mu\)M, 500 \(\mu\)M and 600 \(\mu\)M, respectively. A value of 0.708 ps was calculated for this profile, in the case of sample at 300 \(\mu\)M. Relaxation of this mode is the most stable for dAMP solution, with respect to its concentration. For the band around 1591 cm\(^{-1}\) (C5C4-C4N3), the global relaxation time increases in the interval 200-300 \(\mu\)M, decreases at 400 \(\mu\)M and is equal to 0.708 ps for 500 \(\mu\)M and 600 \(\mu\)M.

### Table 1

**Total Half Bandwidths (cm\(^{-1}\)) of UVRR Vibrational Markers and Global Relaxation Times of Functional Groups, Characterizing dAMP Molecules at Different Concentrations.**

UV resonance Raman spectra were excited with a 275 nm laser line.

| \(\nu_{\text{max}}\) (cm\(^{-1}\)) | \(\Delta\nu_{1/2}\) (cm\(^{-1}\)) | \(\tau_{UVRR}\) (ps) | Tentative assignment [26,27] |
|-----------------|-----------------|-----------------|-----------------|
| **dAMP 200**    |                 |                 |                 |
| 1318            | 19\(\pm\)1      | 0.559\(\pm\)0.029 | N9C8+N3C2+\(\delta\)C8H-\(\delta\)C2H |
| 1433            | 14\(\pm\)1      | 0.758\(\pm\)0.054 | C4N9-\(\delta\)C8H |
| 1494            | 15.5\(\pm\)1    | 0.685\(\pm\)0.044 | -\(\delta\)C2H-N9C8+\(\delta\)C8H |
| 1591            | 14.2\(\pm\)1    | 0.748\(\pm\)0.053 | C5C4-C4N3 |
| **dAMP 300**    |                 |                 |                 |
| 1315            | 16\(\pm\)1      | 0.663\(\pm\)0.041 | N9C8+N3C2+\(\delta\)C8H-\(\delta\)C2H |
| 1432            | 15\(\pm\)1      | 0.708\(\pm\)0.047 | C4N9-\(\delta\)C8H |
| 1490            | 15\(\pm\)1      | 0.708\(\pm\)0.047 | -\(\delta\)C2H-N9C8+\(\delta\)C8H |
| 1590            | 14\(\pm\)1      | 0.758\(\pm\)0.054 | C5C4-C4N3 |
| **dAMP 400**    |                 |                 |                 |
| 1316            | 15\(\pm\)1      | 0.708\(\pm\)0.047 | N9C8+N3C2+\(\delta\)C8H-\(\delta\)C2H |
| 1430            | 14\(\pm\)1      | 0.758\(\pm\)0.054 | C4N9-\(\delta\)C8H |
| 1492            | 15.5\(\pm\)1    | 0.685\(\pm\)0.044 | -\(\delta\)C2H-N9C8+\(\delta\)C8H |
| 1590            | 14.5\(\pm\)1    | 0.732\(\pm\)0.050 | C5C4-C4N3 |
| **dAMP 500**    |                 |                 |                 |
| 740             | 13.5\(\pm\)1    | 0.786\(\pm\)0.058 | -\(\delta\)N7C8N9-9R+\(\delta\)C5N7C8+C4N9C8 |
| 1181            | 21\(\pm\)1      | 0.506\(\pm\)0.024 | N7C8+N1C2; 3C6NH |
| 1315            | 17\(\pm\)1      | 0.624\(\pm\)0.037 | N9C8+N3C2+\(\delta\)C8H-\(\delta\)C2H |
| 1432            | 14.5\(\pm\)1    | 0.732\(\pm\)0.050 | C4N9-\(\delta\)C8H |
| 1491            | 15.5\(\pm\)1    | 0.685\(\pm\)0.044 | -\(\delta\)C2H-N9C8+\(\delta\)C8H |
| 1591            | 15\(\pm\)1      | 0.708\(\pm\)0.047 | C5C4-C4N3 |
| **dAMP 600**    |                 |                 |                 |
| 740             | 13\(\pm\)1      | 0.817\(\pm\)0.063 | -\(\delta\)N7C8N9-9R+\(\delta\)C5N7C8+C4N9C8 |
| 1180            | 18\(\pm\)1      | 0.590\(\pm\)0.033 | N7C8+N1C2; 3C6NH |
| 1317            | 19\(\pm\)1      | 0.559\(\pm\)0.029 | N9C8+N3C2+\(\delta\)C8H-\(\delta\)C2H |
| 1432            | 14\(\pm\)1      | 0.758\(\pm\)0.054 | C4N9-\(\delta\)C8H |
| 1491            | 15.5\(\pm\)1    | 0.685\(\pm\)0.044 | -\(\delta\)C2H-N9C8+\(\delta\)C8H |
| 1591            | 15\(\pm\)1      | 0.708\(\pm\)0.047 | C5C4-C4N3 |

* Low wavenumber side of the band was read; later on the respective profile was symmetrically analyzed. \(\delta\) - bending mode

Experimental fwhm’s of the UVRR bands vs square root of the dAMP concentrations are presented in Figure 2 for different vibrational modes (1318 cm\(^{-1}\), 1433 cm\(^{-1}\), 1494 cm\(^{-1}\) and 1591 cm\(^{-1}\)). It is shown that the linear character is remarkably pronounced in the patterns of the changes in the bandwidths with the square roots of concentrations [20]. This linear relationship indicates that vibrational dephasing is probably the most efficient relaxation mechanism [11].
Fig. 2. Experimental fwhm’s of the UVRR bands vs square root of the dAMP concentration, for the following vibrational modes:
A. 1318 cm⁻¹; B. 1433 cm⁻¹; C. 1494 cm⁻¹; D. 1591 cm⁻¹

Relaxation of the 1430 cm⁻¹ (C4N9-δC8H) oscillator of dAMP in aqueous solutions. Experimental CF and its theoretical modelling

The band near 1430 cm⁻¹ (C4N9-δC8H) in the UV resonance Raman spectrum of dAMP molecule, at the concentration of 400 μM in aqueous solution, shows a fairly symmetric profile without shoulders or splitting (Fig. 1C). This mode allows additional dynamics analysis, namely vibrational band shape analysis through time correlation function (CF) concept. This data treatment is presented in the followings.

The main contribution to the experimental UVRR CF is of vibrational type, due to an adequate selection of the investigated system.

A comparison between natural logarithm of the UV resonance Raman correlation function (CF) for 1430 cm⁻¹ (C4N9-δC8H) band of dAMP, at about 400 μM concentration in aqueous system and the natural logarithm of CFs corresponding to mathematical profiles Gauss and Lorentz (or Cauchy) having the same fwhm’s as the experimental band, is presented in Figure 3. As seen in this figure, the experimental profile fits well with a Gauss one, with Lorentz (or Cauchy) corrections in its wings.

Fig. 3. The comparison between natural logarithm of the UV resonance Raman correlation function (CF) for 1430 cm⁻¹ (C4N9-δC8H) band of dAMP molecule, at about 400 μM concentration in aqueous solution and the natural logarithm of CFs corresponding to mathematical profiles with the same fwhm as the experimental band: black square - experimental CF; dash red line - Gaussian CF; blue line - Lorentzian (or Cauchy) CF
Some models have been developed for vibrational dephasing, as one of the most important contribution to vibrational relaxation, based on Gaussian-Markovian theory [9,14,26-30]. According to Kubo-Rothschild’s model the logarithm of the vibrational correlation function \( G_v(t) \) is expressed by:

\[
\ln G_v(t) = -\left\langle \omega^2 \right\rangle \tau_c^2 \left[ \exp\left(-\frac{t}{\tau_c}\right) - 1 \right] + \tau_c t
\]

(2)

where \( \tau_c \) is the modulation time (time between collisions with the solvent) and \( \left\langle \omega^2 \right\rangle = 4\pi^2 M_2 \) is the mean squared frequency displacement around the band centre.

Furthermore, Oxtoby [30] proposed another equation for this CF:

\[
\ln G_v(t) = -\left\langle \omega^2 \right\rangle \tau_c^2 \ln \cosh\left(-\frac{t}{\tau_c}\right)
\]

(3)

In the followings we will present for band near 1430 cm\(^{-1}\) the theoretical modelling of the natural logarithm of experimental CF with Kubo-Rothschild and Oxtoby models. Due to the fact that there is not a good coincidence between the experimental correlation function (ln CF) and the proposed theoretical models (Kubo-Rothschild and Oxtoby) by employing the experimental second moment, we tried to find the corresponding second band moment for which there is quite a superposition of the experimental correlation function (ln CF) with the proposed models, respectively. So, the value of the modulation amplitude (\( \left\langle \omega^2 \right\rangle \)) was changed until a good coincidence between the natural logarithm of experimental CF and the vibrational dephasing models was obtained. The corresponding modulation time (\( \tau_c \) for Kubo-Rothschild and Oxtoby models, respectively) was obtained in each case. Besides, second band moments, a measure of the kinetic energy of the mode, were calculated with the following relationship: \( M_2 [\text{cm}^2] = 28.33 \times \left\langle \omega^2 \right\rangle [\text{rad. ps}^{-2}] \).

In the next figures this mathematical procedure was applied. Deviations of the theoretical models as compared to the experimental profile, respectively, are to be observed in Figures 4-6. The best fit parameters have been found for Figures 7-8. Also, the spectral and relaxation parameters corresponding to the 1430 cm\(^{-1}\) UVRR vibrational profile for the dAMP (400 μM) sample, are presented in Table 2. Analyzing the values of the modulation speed in this table, a slow modulation regime of the vibrational frequencies is to be observed. Better fits are to be observed upon decreasing of \( \left\langle \omega^2 \right\rangle \). Kubo-Rothschild’s model fits better than Oxtoby’s one the experimental correlation function (ln CF) for values of the modulation amplitude (\( \left\langle \omega^2 \right\rangle \)) ranging between 1.5-1.41 rad. ps\(^{-2}\).

For the modulation amplitude \( \left\langle \omega^2 \right\rangle = 1.40 \text{ rad. ps}^{-2} \), both Kubo-Rothschild and Oxtoby models fit well the experimental correlation function (ln CF).

Fig. 4. The comparison between natural logarithm of the UV resonance Raman correlation function (CF) for 1430 cm\(^{-1}\) (C4N9-6C8H1) band of dAMP molecule, at about 400 μM concentration in aqueous solution (black line) and

A. Kubo-Rothschild’s fitted equation (red line): \( \left\langle \omega^2 \right\rangle = 2 \) and \( \tau_c = 1.35 \);
B. Oxtoby’s fitted equation (red line): \( \left\langle \omega^2 \right\rangle = 2 \) and \( \tau_c = 0.87 \);
C. Kubo-Rothschild’s fitted equation (red line): \( \left\langle \omega^2 \right\rangle = 1.8 \) and \( \tau_c = 1.1 \);
D. Oxtoby’s fitted equation (red line): \( \left\langle \omega^2 \right\rangle = 1.8 \) and \( \tau_c = 1.1 \)
We have found, that upon lowering the kinetic energy stored in this mode, a more efficient vibrational excitation energy transfer from the molecular group to the neighbourhood is obtained. The best fit parameters are connected with the most efficient relaxation mechanism—vibrational dephasing or population relaxation, characterizing biomolecular interactions. The efficiency of this energetic transfer is connected with the biological function of DNA and its components. For example, the respective mode [1430 cm\(^{-1}\) \((C4N9-δC8H)\)], as an energy carrier, might play an important role in molecular recognition or in other interactions of dAMP molecule.
Fig. 7. The comparison between natural logarithm of the UV resonance Raman correlation function (CF) for 1430 cm$^{-1}$ (C4N9-δC8H) band of dAMP molecule, at about 400 μM concentration in aqueous solution (black line) and A. Kubo-Rothschild’s fitted equation (red line): $<\omega^2>=1.45$ and $\tau_c=9.78$; B. Oxtoby’s fitted equation (red line): $<\omega^2>=1.45$ and $\tau_c=2.78$; C. Kubo-Rothschild’s fitted equation (red line): $<\omega^2>=1.42$ and $\tau_c=1.54$; D. Oxtoby’s fitted equation (red line): $<\omega^2>=1.42$ and $\tau_c=3.53$.

Fig. 8. The comparison between natural logarithm of the UV resonance Raman correlation function (CF) for 1430 cm$^{-1}$ (C4N9-δC8H) band of dAMP molecule, at about 400 μM concentration in aqueous solution (black line) and A. Kubo-Rothschild’s fitted equation (red line): $<\omega^2>=1.41$ and $\tau_c=19.02$; B. Oxtoby’s fitted equation (red line): $<\omega^2>=1.41$ and $\tau_c=3.94$; C. Kubo-Rothschild’s fitted equation (red line): $<\omega^2>=1.40$ and $\tau_c=2.50$; D. Oxtoby’s fitted equation (red line): $<\omega^2>=1.40$ and $\tau_c=4.53$.
Table 2
THE SPECTRAL AND RELAXATION PARAMETERS, CORRESPONDING TO THE 1430 cm⁻¹ UVRR VIBRATIONAL PROFILE, FOR THE dAMP (400 μM) SAMPLE: WAVE NUMBER (ωmax); MODULATION AMPLITUDE (<ω²>); SECOND BAND MOMENT (M₂); RELAXATION TIME CALCULATED BY INTEGRATING THE EXPERIMENTAL CF (τUVRR;CF); MODULATION SPEED (<ω²>¹/₂ τc; FOR KUBO-ROTHSCHILD AND OXTOBY MODELS, RESPECTIVELY); MODULATION TIME (τc FOR KUBO-ROTHSCHILD AND OXTOBY MODELS, RESPECTIVELY).

| ωmax (cm⁻¹) | <ω²> (rad. ps⁻²) | M₂ (cm²) | τUVRR;CF (ps) | <ω²>¹/₂ τc; KR | <ω²>¹/₂ τc; Ox | τc;KR (ps) | τc;Ox (ps) |
|-------------|-----------------|----------|---------------|-----------------|-----------------|-----------|----------|
| 1430        | 2               | 56.66    | 0.943         | 1.91            | 1.23            | 0.87      | 1.10     |
| 1430        | 2               | 50.99    | 0.943         | 1.48            | 1.48            | 1.10      | 1.10     |
| 1430        | 1.8             | 48.16    | 0.943         | 2.85            | 1.68            | 1.29      | 1.29     |
| 1430        | 1.6             | 45.33    | 0.943         | 11.98           | 3.52            | 1.58      | 1.58     |
| 1430        | 1.55            | 43.91    | 0.943         | 5.55            | 2.25            | 1.81      | 1.81     |
| 1430        | 1.50            | 42.49    | 0.943         | 7.48            | 6.11            | 2.11      | 2.11     |
| 1430        | 1.50            | 42.49    | 0.943         | 7.48            | 6.11            | 2.11      | 2.11     |
| 1430        | 1.45            | 41.08    | 0.943         | 11.77           | 3.35            | 9.78      | 2.78     |
| 1430        | 1.45            | 41.08    | 0.943         | 11.77           | 3.35            | 9.78      | 2.78     |
| 1430        | 1.42            | 40.23    | 0.943         | 1.84            | 4.21            | 3.53      | 3.53     |
| 1430        | 1.41            | 39.94    | 0.943         | 22.58           | 19.02           | 3.94      | 3.94     |
| 1430        | 1.41            | 39.94    | 0.943         | 22.58           | 19.02           | 3.94      | 3.94     |
| 1430        | 1.40            | 39.66    | 0.943         | 2.95            | 5.36            | 2.50      | 2.50     |
| 1430        | 1.40            | 39.66    | 0.943         | 2.95            | 5.36            | 2.50      | 2.50     |

Conclusions
In this work relaxation processes of functional groups, characterizing 2’-deoxyadenosine-5’-monophosphate (dAMP) compound, were investigated as a function of nucleotide concentration, in aqueous solution (200-600 μM), by UV resonance Raman spectroscopy (UVRR). The full-widths at half-maximum (fwhm’s) studied in this case vary between 13 - 21 cm⁻¹ and the corresponding global relaxation times were found in the range 0.506 - 0.817 ps. Also, the sensitivity of half bandwidths to 2’-deoxyadenosine-5’-monophosphate concentration was analyzed. For aqueous solutions of dAMP molecules, we can suppose that the dominant relaxation mechanism is the vibrational one. From various vibrational relaxation mechanisms, vibrational dephasing seems to be the most important, as indicated from the linear relationship between the true fwhm’s and the square roots of concentrations.

Also, the 1430 cm⁻¹ UVRR band corresponding to ν(C4N9-δC8H) oscillator of dAMP molecule in aqueous solution (400 μM), was selected for band shape analysis. Vibrational dephasing accomplishes an efficient vibrational relaxation of this mode. Theoretical models (Kubo-Rothschild and Oxtoby) for this process have yielded significant relaxation parameters. The value of the modulation amplitude (<ω²>) was varied until a good coincidence between the experimental correlation function (ln CF) and the vibrational dephasing models was obtained, respectively. The best fit parameters are connected with the most efficient relaxation mechanism—vibrational dephasing or population relaxation, characterizing biomolecular interactions. To our knowledge this is the first UVRR study on nucleotide vibrational band shape analysis through time correlation function concept.

This analysis might be used in future studies concerning analogous compounds of DNA and RNA constituents obtained by synthesis, having biological or medical relevance, in particular with an anticancer, antiviral and antibacterial activity.

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References

1. KUMAR PAL, S., PEON, J., ZEWAIL, A.H., Chem. Phys. Lett., 363, 2002, p. 57; Ultrafast decay and hydration dynamics of DNA bases and
mimics.

2. DEVI, T.G., UPADHYAY, G., Spectrochim. Acta A: Molec. and Biomolec. Spectrosc., 91, 2012, p. 106; Solvent dependent Raman bandshape
analysis on C=O containing molecules: Vibrational relaxation study.

3. RABADANOV, K.S., GARUFOV, M.M., ALIEV, A.R., AKHMEDOV, I.R., KAKAGASANOV, M.G., KIRILLOV, S.A., Russ. Metall.+, 8,
2011, p. 760; Vibrational dephasing of the perchlorate ion in LiClO4 melt.

4. RAMAKRISHNAN, V., SARUA, A., KUBALL, M., ABDULLAH, A.F., J. Raman Spectrosc., 41, 2010, p. 320; Raman spectroscopic studies of
vibrational relaxation and non-coincidence enhanced benzaldehyde binary mixtures.

5. GIORGINI, M.G., ARCIONI, A., POLIZZI, C., MUSSO, M., OTTAVIANI, P., J. Chem. Phys., 120, 2004, p. 4969; Relaxation processes of the
liquid crystal ME6N in the isotropic phase studied by Raman scattering experiments.

6. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

7. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

8. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

9. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

10. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

11. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

12. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

13. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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14. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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15. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

16. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

17. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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18. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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19. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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20. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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21. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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22. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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23. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

24. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

25. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

26. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

27. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

28. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

29. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
alkyl benzene solutions.

30. KALAMPOUNIAS, A.G., TSILOMELEKIS, G., BOGHOSIAN, 2, 1981, p. 760; Vibrational relaxation and non-radiative decay processes in
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