Atom-to-atom interactions for atomic layer deposition of trimethylaluminum on Ga-rich GaAs (001)-4 × 6 and As-rich GaAs(001)-2 × 4 surfaces: a synchrotron radiation photoemission study

Tun-Wen Pi1*, Hsiao-Yu Lin2, Ya-Ting Liu3, Tsung-Da Lin3, Gunther K Wertheim4*, Jueinai Kwo2* and Minghwei Hong3*

Abstract

High-resolution synchrotron radiation photoemission was employed to study the effects of atomic-layer-deposited trimethylaluminum (TMA) and water on Ga-rich GaAs(001)-4 × 6 and As-rich GaAs(001)-2 × 4 surfaces. No high charge states were found in either As 3d or Ga 3d core-level spectra before and after the deposition of the precursors. TMA adsorption does not disrupt the GaAs surface structure. For the (4 × 6) surface, the TMA precursor existed in both chemisorbed and physisorbed forms. In the former, TMA has lost a methyl group and is bonded to the As of the As-Ga dimer. Upon water purge, the dimethylaluminum-As group was etched off, allowing the now exposed Ga to bond with oxygen. Water also changed the physisorbed TMA into the As-O-Al(CH3)2 configuration. This configuration was also found in 1 cycle of TMA and water exposure of the (2 × 4) surface, but with a greater strength, accounting for the high interface defect density in the mid-gap region.

Keywords: GaAs, Al2O3, Atomic layer deposition, Synchrotron radiation photoemission

PACS: 73.90.+f, 79.60.Jv

Background

Atomic layer deposition (ALD) facilitates the deposition of a dielectric oxide onto a GaAs surface. The process differs from the one used for the deposition of ALD oxide on Si, where an OH group on the semiconductor is required to initiate the deposition. Bonding of the oxide on the III-V semiconductor is accessible to investigation with high-resolution synchrotron radiation photoemission. It provides unprecedented, precise information about the interfacial electronic structure. This information is vital because the interfacial trap density ($D_{it}$) governs the performance of GaAs-based devices. In order to obtain consistent information, the III-V surface must be free of impurities, such as oxygen, and other defects prior to the ALD process. Only when this condition is satisfied will the true interfacial electronic structure be revealed.

The attempt to prepare a clean GaAs(001) surface has generally been patterned on the procedure used to obtain a clean Si(001) surface. That neglects the fundamental difference between the surface properties and reconstruction of a III-V semiconductor and an elemental one like Si. The reconstructed Si(001)-2 × 1 surface consists of rows of buckled dimers, with charge transfer between the tilted atoms, and is rich in dangling bonds that trap impurities. Surface pretreatment is required prior to a final anneal in an ultra-high-vacuum end station prior to synchrotron radiation photoemission (SRPES) measurements. The pre-treatment due to Ishizaka and Shiraki [1] has come into general use. It leaves a thin oxide film on a clean Si surface that is readily removed by annealing in vacuum [2,3]. The effectiveness of this procedure has been demonstrated in [2], which shows the analysis of 2p core-level data from a clean reconstructed Si(001) surface. The photoemission...
spectra from the first three surface layers labeled S(0), S(1), and S(2) are identified and have intensities consistent with the expected escape depth.

For multi-element (In)GaAs, a common method of surface pretreatment prior to in-vacuum annealing is As capping [4] by thermal annealing in As$_2$ flux [5], followed by a chemical rinse [6]. Subsequent in-vacuum annealing of these samples removes the more volatile As and produces an oxygen-free surface, but one that does not have the desired surface Ga/As ratio. It turns out to be low, say 0.73 [4], compared with an untreated sample, say 1.26 (not shown). The process has clearly changed the surface stoichiometry as well as the surface reconstruction. To be specific, ALD of Al$_2$O$_3$ with trimethylaluminum (TMA) and water on the treated GaAs(001) with ammonia or ozone often left As-As dimers at the interface, resulting in significant frequency dispersion in the C-V characteristic curve [7-9]. This conventional cleaning process does not reproduce the clean reconstructed surface and must be adjudged a failure. The resulting uncertainty regarding the chemistry and reconstruction of the surface prevents an understanding of the nature of the interaction with adsorbates and stands in the way of systematic improvements. It impacts both work on the interfacial electronic structure of high-κ dielectric oxides/(In)GaAs [10-12] and spintronics based on Fe$_3$Si/GaAs [13,14].

In this work, we present a high-resolution core-level SRPES investigation of the electronic structure of the clean, Ga-rich GaAs(001)-4 × 6 surface, followed by the characterization of the surface after 1 cycle of ALD of, first, TMA and then water H$_2$O onto the TMA-covered surface. For comparison, we also present the data of 1 cycle of TMA and H$_2$O on As-rich GaAs(001)-2 × 4. We note that the ALD precursors were exposed onto a surface with a long-range order, a condition of that has not been previously achieved in work with GaAs.

**Method**

The samples were fabricated in a multi-chamber growth/analysis system, which includes a GaAs-based molecular beam epitaxy (MBE) chamber, an ALD reactor, and many other functional chambers [15,16]. These chambers are connected via transfer modules, which maintain ultra-high vacuum of 10$^{-10}$ Torr. Thus, pristine surfaces were obtained during the sample transfer. MBE was employed to grow Si-doped GaAs (1 to 5 × 10$^{17}$ cm$^{-3}$) onto 2-in. n-GaAs(100) wafers. ALD was employed to high κ dielectrics on freshly MBE-grown GaAs. The samples were transferred in vacuo into a portable module kept at 2 × 10$^{-10}$ Torr and transported to the National Synchrotron Radiation Research Center located in Taiwan for SRPES measurements. Photoelectrons were collected with a 150-mm hemispherical analyzer (SPECS, Berlin, Germany) in an ultra-high vacuum chamber with a base pressure of approximately 2 × 10$^{-10}$ Torr. The overall instrumental resolution was better than 60 meV, and the binding energy was established in accordance with the Fermi edge of Ag.

**Results and discussion**

The surface reconstruction of GaAs(001) was first checked with reflection high-energy electron diffraction in the molecular beam epitaxial growth chamber and then verified with low-energy electron diffraction (LEED) in the photoemission chamber. The LEED pattern is shown in Figure 1a. It consists of sharp 4 × 6 spots and third-order
streaks along the [110] direction. The streaking pattern indicates that the surface contains small domains of (6 × 6) or c(8 × 2) reconstruction. The low background intensity indicates that the surface is smooth with a great long-range order. Recently, Ohtake et al. proposed that for the GaAs(001)-4 × 6 surface, the unit cell is composed of four As atoms in the faulted position, two tilted As-Ga dimers and two Ga-Ga dimers in the subsurface layer [17]. The atomic structure of the Ohtake model is shown in Figure 1b.

Figure 1c displays the As 3d and Ga 3d core-level spectra of a clean Ga-rich n-GaAs(001)-4 × 6 surface taken in various angles from the normal emission to 60° off-normal emission. The excitation photon energies were set at 85 and 65 eV for As and Ga states, respectively. The estimated escape depth is approximately 0.3 to 0.5 nm. A visual inspection of the As (Ga) 3d photoemission data identifies a feature bulged out at low (high) binding energy, suggesting that the line shape contains components in addition to the main bulk line. In fact, deconvolution of the As 3d core-level spectrum shows four components. Accordingly, we set up a model function with four spin-orbit pairs as well as a power-law background and a plasmon- or gap-excitation-energy loss tail. The background and loss tail are represented by least squares adjustable parameters that are included in the model function. The background is represented by four parameters: a constant, a slope, and a power-law that is quite successful in representing the degraded electrons from shallower levels. In the energy range of the 3d spectra, the loss tail is almost entirely due to electron-hole pair excitations in the semiconductor. In GaAs, there are none that are smaller than the 1.42-eV bandgap, which implies that almost all of the line structure remains unaffected by the loss tail. Background subtraction prior to fitting meets with a fundamental objection. It destroys the statistical relationship between the number of counts in the data point and its uncertainty, preventing $\chi^2$ from reaching unity for a perfect fit and interfering with the assessment of the quality of the fit.

The fact that the resolved components in the deconvolute exhibit nearly equal widths suggests that the lifetime is the same for all components. The residual differences in width are presumably due mainly to small differences in the phonon or inhomogeneous broadening of bulk and surface components. It is worth noting that a reliable least

Figure 2 Analysis of the core-level spectra for the clean Ga-rich GaAs(001)-4 × 6 surface. (a) As 3d state, $\theta_e = 0^\circ$, (b) As 3d state, $\theta_e = 60^\circ$, (c) Ga 3d state, $\theta_e = 0^\circ$, and (d) Ga 3d state, $\theta_e = 60^\circ$. 
squares adjustment is readily obtained provided the model function has a multi-parameter global minimum. A multitude of unconstrained width parameters tend to produce local minima defining erroneous, unphysical parameter values. The width parameters were accordingly constrained as needed.

The representative fit to the As and Ga 3d states of the clean GaAs(001)-4 × 6 surface are shown in Figure 2. The numerical results were obtained using the nonlinear Marquardt algorithm. The surface core-level shifts (SCLSs) of the Ga 3d state for the S1', S2', and S3' components relative to the bulk at 19.58 eV are −0.302, +0.251, and +0.613 eV, respectively. The Gaussian widths of the bulk and surface are 0.33 and 0.45 eV, respectively. For the As 3d state, the S1, S2, and S3 components relative to the bulk located at 40.43 eV (the 3d_{5/2} state) were found to be +0.159, −0.249, and −0.599 eV, respectively. A ‘+’ or ‘−’ sign indicates a shift towards a higher or lower binding energy, respectively. The Gaussian width is about 0.31 eV. The lifetime is 0.22 eV. In Figure 2b,d, the change in intensity of the components at 60° emission angle is displayed, clearly identifying the surface components. The smallest As component, S3, is most likely associated with the As in the tilted As-Ga dimers in the defaulted terrace. The shifted magnitude of component S3 is the greatest among those reported in the literature, suggesting that the tilted angle of the dimer is great so as to cause a large charge transfer.

Figure 3 displays a fit to the TMA-exposed surface prior to exposure to H₂O. As shown in Figure 3a, two Al 2p states are well resolved with an energy separation of 0.650 eV. The one with lower binding energy is associated with a charge transfer from As to Al. This is possible when a methyl ligand is replaced by a direct bond to an As atom. Considering that the GaAs(001)-4 × 6 surface is ‘As-terminated’ and component S3 shows a negative SCLS, we assumed that dimethylaluminum (DMA) bonds with the dangling bond of the As in the As-Ga dimer.

Because the high-binding-energy Al 2p state remains in the same position and with similar line width after the subsequent water purge, the TMA precursor must have maintained the Al in the molecular charge state while residing on the surface. That indicates that this TMA does not form a bond with a surface atom. That is in agreement with the absence of a new surface As level and leads to the conclusion that the TMA is physisorbed on the S1 As atoms.

For the As 3d core-level spectrum, the TMA-exposed surface reveals only minor changes from the clean surface. First, the widths of both top-surface S1 and S3 components are 15% to 20% broader than the subsurface S2 component. Second, the SCLS of the S1 component becomes 0.056 eV without changing the strength. Third, the intensity of the S3 component slightly decreases concurrently with a slight increase of the S2 intensity. Because the Al in DMA bonds with S3 As atom, this As underneath the Al behaves as a subsurface atom. Consequently, the energy position coincides with that of the S2 state, thereby increasing the S2 intensity slightly.
Figure 3c presents a fit to the Ga 3d core-level spectrum. The Ga 3d states remain virtually unaltered, indicating that the TMA precursor has not disturbed the Ga layer.

Figure 4 displays a fit to the spectra after 1 cycle of TMA and H₂O purges. The Al 2p state now exists as a single peak without any sign of the component identified with DMA. This suggests that the H₂O precursor has etched off the attached Al-(CH₃)₂ species that bonded to the As in the As-Ga dimer. Removal of the As atoms exposes the previously dimerized Ga atom which now becomes oxidized as shown in Figure 4c, where the oxidized Ga* state appears with SCLS of +0.892 eV. Note that the area of the S2 state retains the magnitude in the clean surface.

Figure 4b exhibits As-induced states labeled as As* with SCLSs of +0.680 eV. The energy separation of the As* and S1 states is 0.432 eV, which remains constant in the greater cycles of deposition (not shown), indicating that the As* state originated from the S1 As atoms. Because the SCLS of the As* state becomes more positive than that of the S1 state, under the influence of water, the adsorbed TMA precursor must undergo a change of bonding configuration to become a charge acceptor for the affiliated As atom. Because no similar Al-X state...
appears in the Al 2p core-level spectrum, water then affects the TMA molecule that is physisorbed on As in a way that allows the interfacial S1 As to become an As-O-Al configuration, where the surface is further terminated with a hydroxyl group.

Figure 5a shows a fit to the As 3d core-level spectrum for the clean As-rich GaAs(001)-2 × 4 surface. The β2(2 × 4) model is commonly believed to represent the surface reconstruction, where the top surface layer is characterized as two rows of As-As dimers separated by itself from an As-As dimer located in the third layer. As can be seen in Figure 5a, three surface components were resolved. With reference to an off-normal spectrum (not shown), both the S1 and S3 components are identified with the surface As-As dimers because of the intensity enhancement. In fact, components S3 and S1 are associated with the clean As-rich GaAs(001)-2 × 4 surface. The S3 component has been re-termin. In fact, components S3 and S1 are associated with surface As-As dimers because of the intensity enhancement. The model is commonly believed to represent the surface reconstruction, where the top surface layer is characterized as two rows of As-As dimers separated by itself from an As-As dimer located in the third layer. As can be seen in Figure 5a, three surface components were resolved. With reference to an off-normal spectrum (not shown), both the S1 and S3 components are identified with the surface As-As dimers because of the intensity enhancement.

Conclusions
We have presented a microscopic view of in situ atomic layer deposition of TMA and H2O precursors on atomically clean GaAs(001) surfaces in both 4 × 6 and 2 × 4 reconstructions. For the Ga-rich 4 × 6 surface, the precursors partially and selectively bond with the surface atoms without disturbing the atoms in the subsurface layer. TMA is dissociative on As in the As-Ga dimer but is physisorbed on As that is threefold Ga-coordinated. Water drastically alters the TMA-covered surface, etching off the DMA along with its As, resulting in Ga-O bonding for the subsequent deposition of Al2O3. At the same time, it transforms the configuration of the physisorbed TMA to bond strongly with As. On the As-rich 2 × 4 surface, 1 cycle of TMA and H2O entirely passivates the surface As dimer bonds.

Abbreviations
ALD: Atomic layer deposition; D0: Interfacial trap density; DMA: Dimethylaluminum; LEED: Low-energy electron diffraction; SCLS: Surface core-level shift; SRPES: Synchrotron radiation photoemission; TMA: Trimethylaluminum.

Competing interests
The authors declare that they have no competing interests.

Authors’ contributions
All authors have contributed to the final manuscript of the present work. MH, JK, and TWP defined the research topic. TDL provided the GaAs sample. YTL prepared the precursor-purged interfaces. HYL acquired the photoemission data. TWP wrote the paper. GKW and MH provided critical comments on the draft manuscript. All authors read and approved the final manuscript.

Authors’ information
TWP is a research scientist at the National Synchrotron Radiation Center, Hsinchu, Taiwan. GKW is the President of Woodland Consulting in the USA. JK is a full professor in the Department of Physics, National Tsinghua University, Hsinchu, Taiwan. MH is a full professor in the Graduate Institute of Applied Physics and Department of Physics, National Taiwan University, Taipei, Taiwan. TDL is a postdoctoral fellow in the Graduate Institute of Applied Physics and Department of Physics, National Taiwan University, Taipei, Taiwan. HYL is a graduate student in the Department of Physics, National Tsinghua University, Hsinchu, Taiwan. YTL is a graduate student in the Department of Materials Science, National Tsinghua University, Hsinchu, Taiwan.

Acknowledgments
This project was supported by the National Nano Projects (NSC-98-2120-M-007-002) and NSC 99-2112-M-213-004-MY3 of the National Science Council in Taiwan. We also acknowledge the support of AOARD.

Received: 6 November 2012 Accepted: 7 March 2013
Published: 12 April 2013

References
1. Ishizaka A, Shiraki Y: Low temperature surface cleaning of silicon and its application to silicon MBE. J Electrochem Soc 1986, 133(4):666–671.
2. Pi TW, Cheng CP, Wertheim GC: The reaction of Si(001)-2×1 with magnesium and calcium. J Appl Phys 2011, 109:033703.
3. Pi TW, Hong IH, Cheng CP, Wertheim GC: Surface photoemission from Si(100) and inelastic electron mean-free-path in silicon. J El Spectr Rel Phen 2000, 107:163–176.
4. Thompson JDW, Neal JR, Shen TH, Morton SA, Tobin JG, Dan Widdall G, Matthew JAD, Greig D, Hopkinson M: The evolution of Ga and As core levels in the formation of FeGaAs(001): a high resolution soft x-ray photoelectron spectroscopic study. J Appl Phys 2008, 104:024516.
5. Laukkanen P, Perala RE, Vaira R-L, Väyynen U, Kuzmin M, Saduwski J: Electronic and structural analysis of Si-2-8-As(100)-2×4 and (2×8) surfaces. Phys Rev B 2004, 69:205323.
6. Tereshchenko OE: Preparation of As-rich (2x4)-III-As (001) surfaces by wet chemical treatment and vacuum annealing. Phys Stat Sol C 2010, 7:264.
7. Hinkel CL, Sonnet AM, Vogel EM, McDonnell S, Hughes GJ, Milojevic V, Lee S, Aguero-Tostado FS, Choi KJ, Kim J, Wallace RM: Frequency dispersion reduction and bond conversion on n-type GaAs by in situ surface oxide removal and passivation. Appl Phys Lett 2007, 91:163512.
8. Shahrijerdi D, Akyol T, Ramon M, Garcia-Gutiérrez DT, Tutuc E, Banerjee SK: Self-aligned inversion-type enhancement-mode GaAs metal-oxide-semiconductor field-effect transistor with Al2O3 gate dielectric. Appl Phys Lett 2008, 92:203505.
9. Hinkel CL, Milojevic V, Vogel EM, Wallace RM: Surface passivation and implications on high mobility channel performance. Microelectron Eng 2009, 86:1544–1549.
10. Hong MW, Kwo JR, Tsai PC, Chang YC, Huang ML, Chen CP, Lin TD: III-V metal-oxide-semiconductor field-effect transistors with high k dielectrics. Jpn J Appl Phys 2007, 46(5B):3167–3180.
11. Robertson J, Lin L: Bonding principles of passivation mechanism at III-V oxide interfaces. Appl Phys Lett 2011, 99:222906.
12. Chang YH, Lin CA, Liu Y, Chang TH, Lin HY, Huang ML, Lin TD, Pi TW, Kwo J, Hong M: Effective passivation of InGaAs by H2O surpassing Al2O3 via in-situ atomic layer deposition. Appl Phys Lett 2012, 101:72104.
13. Hong M, Chen HS, Kwo J, Kortan AR, Mannaerts JP, Weir BE, Feldman LC: MBE growth and properties of Fe3(Al, Si) on GaAs(100). J Crystal Growth 1991, 111:984–988.
14. Ionescu A, Kaz CAF, Trypiniotis T, Gürtler CM, Garcia-Miquel H, Bland JAC, Vickers ME, Dalgliesh RM, Langridge S, Bugoslavsky Y, Miyoshi Y, Cohen LF, Diebold KRA: Structural, magnetic, electronic, and spin transport properties of epitaxial Fe3Si/GaAs(001). Phys Rev B 2005, 71:094401.
15. Hong M, Mannaerts JP, Bowers JE, Kwo J, Passlack M, Hwang WY, Tu LW: Novel Ga$_2$O$_3$(Gd$_2$O$_3$) passivation techniques to produce low $D_{it}$ oxide-GaAs interfaces. J Crystal Growth 1997, 175/176:422–427.

16. Chang YH, Huang ML, Chang P, Lin CA, Chu YJ, Chen BL, Hsu CL, Kwo J, Pi TW, Hong M: Electrical properties and interfacial chemical environments of in-situ atomic layer deposited Al$_2$O$_3$ on freshly molecular beam epitaxy grown GaAs. Microelectron Eng 2011, 88:440–443.

17. Ohtake A, Kocan P, Seino K, Schmidt WG, Koguchi N: Ga-rich limit of surface reconstructions on GaAs(001): atomic structure of the (4x6) phase. Phys Rev Lett 2004, 93:266101.

18. Chang YC, Merckling C, Penaud J, Lu CY, Wang WE, Dekoster J, Meuris M, Caymax M, Heyns M, Kwo J, Hong M: Effective reduction of interfacial traps in Al$_2$O$_3$/GaAs (001) gate stacks using surface engineering and thermal annealing. Appl Phys Lett 2010, 97:12901.

19. Chang YC, Chang WH, Merckling C, Kwo J, Hong M: Inversion-channel GaAs(100) metal-oxide-semiconductor field-effect-transistors using molecular beam deposited Al$_2$O$_3$ as a gate dielectric on different reconstructed surfaces. Appl Phys Lett 2013, 102:093506.