A Part Per Trillion Isotope Ratio Analysis of 90Sr/88Sr using Energy-Filtered Thermal ionization mass spectrometry

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Abstract

Strontium-90 is a major radioactive nuclide released by nuclear accidents and discharge waste. Input of such radioactive nuclide into earth surface environment causes potential threat of long-term internal exposure when taken up by organism. Rapid and precise measurement of $^{90}$Sr in variety of environmental sample is important to understand the distribution and dynamics of $^{90}$Sr in the local environment after the accident and to assess the effect of radioactive nuclide inputs on bodies. However, previous $^{90}$Sr measurement techniques have drawbacks such as long measurement times for radiometry and high detection limits for mass spectrometry. Here we present a technique to accurately measure a significantly small amount of $^{90}$Sr in natural environmental samples using an energy-filtered thermal ionization mass spectrometry. Our technique achieved a $^{90}$Sr detection limit of 0.23 ag, which corresponds to a $^{90}$Sr activity of 1.2 µBq. The detection limit was lowered by two orders of magnitude compared with the previous mass spectrometric $^{90}$Sr analyses. The ability of our technique will expand the applicability of mass spectrometric $^{90}$Sr survey not only to the rapid $^{90}$Sr survey upon nuclear accidents but also to study a long-term environmental diffusion of radioactive materials using size-limited environmental and biological samples.

Introduction

The moment of history or dynamics of human and environment are precisely memorized as isotope ratios in variety of materials. Understanding the isotope ratio stem from radioactive isotope can realize an aspect of past condition and transition. Among the isotope ratio, that of strontium (Sr) reflects past aspect and/or a fact basing on radioactive decay or mass discrimination effect, and it widely serves indictors as age determination in geological sciences, environmental dynamics, forensic sciences, authenticity judgment, and determination of geographic origin. In particular, the artificial nuclide, radioactive strontium ($^{90}$Sr) records as the memory of nuclear disaster in materials such as Chernobyl or Fukushima Nuclear Power Plants Accidents.

$^{90}$Sr is a pure β-ray emitting radionuclide (0.55 MeV, 100% yield) and a typical fission product of $^{235}$U (5.8%) and $^{239}$Pu (2.1%). Anthropogenic $^{90}$Sr is widely spread over earth surface environments with low background activities because of the use and testing of atomic bombs. Nuclear accidents are another source of anthropogenic $^{90}$Sr emissions into the environment, resulting in local fallout with high activities. $^{90}$Sr was one of the major radioactive nuclides released in the Chernobyl and Fukushima nuclear reactor accidents. When taken up by animals, including humans, $^{90}$Sr has the potential to cause long-term internal exposure. The slight trace of $^{90}$Sr is precisely memorized as an isotope ratio between stable isotope of Sr in variety of samples. Therefore, a precise measurement the isotope ratio of $^{90}$Sr and stable Sr is important to determine the understanding of right dynamics of $^{90}$Sr in the local environment surrounding the accident site.
Conventionally, radiometric methods using solid/liquid scintillators or gas ionization detectors were used to detect $^{90}$Sr in various environmental samples. The radiometric method is fundamentally a counting experiment of the decaying atoms. Because $^{90}$Sr has a half-life of 28.79 years, the fraction of the decaying atom is significantly small: approximately $1.3 \times 10^9$ atoms of $^{90}$Sr corresponds to 1 Bq (or decay per minute) of activity. Therefore, for precise radiometric $^{90}$Sr determination, a relatively large sample size or long analysis time is required. Because $^{90}$Sr is a pure $\beta$ emitter, radiometric determination of $^{90}$Sr includes complex radiochemical procedures and typically requires 2 weeks or more measurement time. This slow sample processing speed is insufficient for an urgent environmental survey in response to a nuclear emergency.

An alternative and less time-consuming method to analyze $^{90}$Sr is mass spectrometry. Inductively coupled plasma mass spectrometry (ICP-MS) is a new mass spectrometric technique for $^{90}$Sr detection. Unlike radiometry, the target of mass spectrometric measurement is all the existing $^{90}$Sr atoms in the sample. The large difference in target numbers suggests that mass spectrometric methods are more sensitive for detecting $^{90}$Sr. However, measuring a minor isotope is a challenging task in mass spectrometry. Mass spectrometry's sensitivity to detect trace $^{90}$Sr is hindered by high noise signals due to isobaric interferences of $^{90}$Zr and molecular ions, as well as peak tailing of the highly abundant $^{88}$Sr ions. ICP-MS's major source of noise signals is $^{90}$Zr. In environmental samples, the ratio between minor $^{90}$Sr and more abundant stable isotopes of Sr or isobaric isotope $^{90}$Zr is extremely high (i.e., at least nine orders of magnitude). ICP ion source efficiently ionizes $^{90}$Zr, which remains in the sample solution in trace amounts even after Sr extraction chemistry. The reaction of the ions with O$_2$ gas is used in a dynamic reaction cell (DRC) technique to reduce the $^{90}$Zr ion transmission. Typically, this technique achieves abundance sensitivities, defined as the intensity ratio between $^{88}$Sr peak tail on m/z = 89.908 and $^{88}$Sr, on the order of $10^{-9}$. Furthermore, the introduction of a new technique, triple quadrupole ICP-MS (ICP-MS/MS), combined with the O$_2$ gas reaction, effectively reduced the $^{90}$Zr ion transmission in the mass spectrometer and lowered the noise signal to 0.1 cps, which correspond to a $^{90}$Sr detection limit of 0.11 fg (0.6 mBq). With a sample size of 4 μg of Sr, the abundance sensitivity for the $^{90}$Sr/$^{88}$Sr ratio achieved by this technique was $5 \times 10^{-12}$. However, reducing the $^{90}$Sr detection limit with the ICP ion source is difficult. The highly efficient ICP ion source mainly produces abundant polyatomic or polyvalent ions from the sample solution's solvents, Ar gas, and trace impurities in the sample solution, and such ions exist across the entire m/z range. It also emits Ar gas-related ions such as Ar$^+$ and ArO$^+$ at significantly high intensities, which may cause non-spectrum interference of the peak tails. All such interferences are increasing the noise level and limiting the sensitivity of ICP-MS analysis. Even with the ICP-MS/MS technique, the background signal was reported as 0.1–0.2 cps when aspirating a blank solution, showing that noise signals of sub-cps level are inevitable while using the ICP ion source.

Thermal ionization mass spectrometry (TIMS) is a standard technique for measuring isotope ratios of Sr, such as $^{87}$Sr/$^{86}$Sr and $^{88}$Sr/$^{86}$Sr. Compared with the ICP ion source, the thermal ionization ion source
is energy-limited, and thus ionization is mostly limited to target elements with very few polyatomic ions. In TIMS, Ar-gas-related species, solvent-related species, and polyvalent ions are mostly absent. With less spectrum and non-spectrum interferences, TIMS can reduce noise signal levels on $^{90}$Sr and thus increase sensitivity upon $^{90}$Sr detection. Recently, several attempts have been made to detect $^{90}$Sr using the TIMS technique. Previous studies have failed to detect traces of $^{90}$Sr in environmental samples.$^{27,28}$ Kavasi et al. measured $^{90}$Sr/$^{88}$Sr ratios of $^{90}$Sr-containing reference materials of wild berry and lake sediment using a sample size of 1000 ng of Sr and reports an abundance sensitivity for the $^{90}$Sr/$^{88}$Sr ratio as $2.1 \times 10^{-10}$, corresponding to a $^{90}$Sr noise signal of 0.77 cps.$^{37}$ However, when the $^{90}$Sr/$^{88}$Sr ratio is lower, their measured $^{90}$Sr/$^{88}$Sr ratio shows a systematic bias toward higher values. To account for such a significant bias, additional empirical “relative bias” errors must be introduced into the analytical uncertainty. This inaccuracy is obvious when the $^{90}$Sr/$^{88}$Sr ratio is lower than $1.2 \times 10^{-9}$ and is likely to be caused by an inaccurate noise correction scheme for the $^{90}$Sr signal. Ito et al. focused on analyzing small-sized samples and constructed an isotope-dilution total-evaporation (ID-TE-) TIMS technique to analyze $^{90}$Sr with a sample size of 5–20 ng of Sr.$^{38}$ Their study did not report abundance sensitivity, but it can be estimated as $2 \times 10^{-8}$ based on the reported analytical conditions of $^{90}$Sr noise level of approximately 5 cps and $^{88}$Sr target intensity of 4 V.$^{38}$ During the TIMS analysis, the ionization of $^{90}$Zr$^+$ is suppressed, whereas Sr is still present and ionizing from the filament$^{38}$ because of the difference in the evaporation and ionization potentials between the two elements. Therefore, the peak tailing of the highly abundant $^{88}$Sr ions is the main source of background signal for $^{90}$Sr in TIMS measurements. Peak tail ions can be efficiently reduced using an energy filtering device placed immediately before the detector. However, none of the previous $^{90}$Sr analyses used an effective energy filtering device to account for the $^{88}$Sr tailing.$^{35–38}$ Kavasi et al. used a WARP energy filter. However, they did note that the WARP energy filter eliminates low energy ions but has no effect on high-mass side peak tailing, implying that the WARP energy filter does not work for $^{88}$Sr peak tail on $^{90}$Sr. The ability of TIMS to detect a significantly small amount of $^{90}$Sr remains unknown.

In this study, we focused on bringing out the performance of modern TIMS instruments to perform an accurate $^{90}$Sr/$^{88}$Sr measurement of environmental and biological samples with low $^{90}$Sr activity using an effective energy filtering technique for $^{90}$Sr detection. The Retarding Potential quadrupole (RPQ) lens act as high selectivity filter for ions with disturbed energy or angle.$^{39}$ The use of RPQ lenses for energy filtering coupled with reduction of the multiple noise signal sources resulted in a significantly low and highly stable noise signal for $^{90}$Sr compared with the previous studies. With significantly low and stable noise signals, an appropriate noise correction scheme was used in this study to allow accurate measurement of $^{90}$Sr/$^{88}$Sr ratio down to $10^{-11}$ level. The measurement of $^{90}$Sr-containing reference materials (plants), as well as environmental and biological samples, demonstrates the ability of our technique.

Results And Discussion
Noise reduction schemes upon $^{90}\text{Sr}$ detection by energy-filtered TIMS. Peak tailing of $^{88}\text{Sr}$, which has the largest abundance among Sr isotopes, is the main limiting factor in TIMS detection of $^{90}\text{Sr}$. The RPQ lenses used in this study eliminate peak tailing from the high-mass and low-mass sides of the $^{88}\text{Sr}$ peak (Figure S1 in supporting information). The RPQ lens parameters must be fine-tuned to effectively eliminate the $^{88}\text{Sr}$ peak tailing on $^{90}\text{Sr}$ while sacrificing only a small amount of ion transmission rate and peak shape. Finally, after the fine-tuning of the RPQ lenses, a 1000-fold reduction of the $^{88}\text{Sr}$ peak tail signal was achieved (Figure 1).

Because the $^{88}\text{Sr}$ peak tail signal has been reduced to approximately 0.01 cps on $^{90}\text{Sr}$, signals from other noise sources must also be controlled at this noise level. The filament material, rather than the sample-derived Zr, is the major source of the isobaric $^{90}\text{Zr}$ signal in TIMS measurement. Typically, filaments were baked out at 4500 mA (approximately 2000°C) before the use. With this level of bake out, about 1 cps and about 0.1 cps of the $^{90}\text{Zr}^+$ signals were observed from Re and Re-ZR filaments, respectively, during Sr isotope ratio measurement even with the presence of Sr sample on the filament. These were not acceptable levels of noise in this study. Therefore, filaments were baked at higher temperatures at a 5.5 A filament current (approximately 2200°C) to reduce the filament-derived Zr$^+$ signal. After this high temperature bakeout, the Re and the Re-ZR filaments tend to show $^{90}\text{Zr}^+$ signals on the 0.01 cps order when heated up to 1550°C without the sample material. Even after this high temperature bakeout, a few Re and Re-ZR filaments still demonstrate higher $^{90}\text{Zr}^+$ emission. Therefore, after the bakeout, all the filaments were inspected for $^{90}\text{Zr}$ intensity to ensure that filaments with higher $^{90}\text{Zr}^+$ emission were not used. The $^{90}\text{Zr}$ intensity threshold at 1550°C was empirically set to 0.04 cps. Note that the actual $^{90}\text{Zr}$ intensity during Sr isotope measurement is negligible (far lower than 0.04 cps) because Zr$^+$ emission is suppressed during the presence of Sr on the filament.

Ito et al. observed that organic material with m/z 90.0 is a significant isobaric interfering molecule on $^{90}\text{Sr}$, and this could be perfectly mass separated by slightly shifting the axial m/z from 89.908 to 89.777. We detected another isobaric interference molecule spectrum at m/z 89.908 with intensities on the order of 0.1 cps or less (Figure S2 in supporting information). The intensity of this noise signal is relatively high at the beginning of the measurement and decreases with time (Figure S3 in Supporting Information), indicating that this noise signal is related to organics or volatiles. The most plausible molecule for this mass spectrum is $^{88}\text{SrH}_2^+$, with hydrogen derived from organics or residual H$_2$O. This noise signal was eliminated by preheating of the sample filament under vacuum before the measurement and by avoiding the use of parafilm/catheter during sample loading.

**Abundance sensitivity and detection limit of the $^{90}\text{Sr}/^{88}\text{Sr}$ measurement.** The $^{90}\text{Sr}$-free NIST SRM-987 was measured using all these noise reduction schemes to validate the noise level and noise stability on
The results of the three analytical sessions with slightly different analytical conditions are summarized in Table 1, Figure 2, and Figure S4 in Supporting Information.

### Table 1 Results of replicate measurements of NIST SRM 987 in three different mass spectrometric sessions

| session     | reference date       | $^{88}\text{Sr}$ [V] | $^{90}\text{Sr}$ \(^{1,2}\) | average darknoise \([\text{cps}]\) | $^{90}\text{Sr}/^{88}\text{Sr}$ \(^{1,2}\) (abundance sensitivity) | 3SD (detection limit) | number of analysis |
|-------------|----------------------|-----------------------|-----------------------------|---------------------------------|------------------------------------------------------------------|----------------------|------------------|
| Bremen      | January 22nd, 2020   | 20.3                  | 0.0118 \((29)\)             | 0.0070                           | 9.3 (3.0) \(\times 10^{-12}\)                                      | 3.7 \(\times 10^{-12}\) | 4                |
| Fukushima-1 | April 3rd, 2021      | 27.2                  | 0.0168 \((34)\)             | 0.0125                           | 1.0 (1.8) \(\times 10^{-11}\)                                      | 3.1 \(\times 10^{-12}\) | 13               |
| Fukushima-2 | June 5th, 2021       | 26.3                  | 0.0136 \((29)\)             | 0.0096                           | 8.3 (1.8) \(\times 10^{-12}\)                                      | 2.7 \(\times 10^{-12}\) | 17               |

1) Errors in the parenthesis are 2SD.

2) These data are not corrected for SEM dark noise.

The average intensity of the $^{90}\text{Sr}$ noise signal of the three analytical sessions were 0.0118 ± 0.0029 (2SD, \(n = 4\)), 0.0168 ± 0.0034 (2SD, \(n = 13\)) and 0.0136 ± 0.0029 (2SD, \(n = 17\)) for Bremen, Fukushima-1st and Fukushima-2nd sessions, respectively. Note that these noise signals include dark noise counts because they were not corrected. Slightly higher noise intensity of Fukushima-1st compared with the Bremen and Fukushima-2nd reflects the fact that higher SEM operating voltage resulted in higher dark noise level in Fukushima-1st session (Table 1). The average dark noise count during the Fukushima-2nd session was 0.0096 cps, accounting for 70% of the total noise signal. This indicates that our noise reduction scheme has successfully reduced noise signals other than the detector’s intrinsic dark noise to the level of 0.004 cps (Figure 2). The small variability of the noise signal indicates that the noise signals are well controlled at this level. In this study, the level and stability of the noise signal achieved are superior compared with the noise signals reported by non-energy filtered TIMS technique\(^{37}\) (0.77 ± 0.82 cps, 2SD) and by ICP-MS/MS technique\(^{28}\) (0.1 cps). Noted that measurement with such a low signal requires a long measurement time (1 h in this study) to count the substantial number of ions. The stable nature of the TIMS ion source allows long measurement with stable analytical conditions.
The abundance sensitivity and detection limit of the \(^{90}\text{Sr}/^{88}\text{Sr}\) ratio can be determined from these data. Finally, the abundance sensitivity \(^{90}\text{Sr}/^{88}\text{Sr}\) ratio achieved in the *Fukushima-2nd* session was \(8.3 \pm 1.8 \times 10^{-12}\) (2SD, \(n = 17\); Table 1). A detection limit of a signal is defined by \(3\sigma\) of the variability of the zero-point (or blank) analysis. Therefore, the detection limit of \(^{90}\text{Sr}/^{88}\text{Sr}\) ratio finally achieved in the *Fukushima-2nd* session was estimated to be \(2.7 \times 10^{-12}\) (Table 1 and Figure S4(c) in Supporting Information). The fact that the detection limit is lower than the abundance sensitivity demonstrates TIMS’s excellent noise stability and its good control on \(^{90}\text{Sr}\).

**Noise correction scheme for \(^{90}\text{Sr}\) signal.** The stable nature of the noise signal allows precise noise correction on \(^{90}\text{Sr}\) measurement. Our noise correction scheme is as follows. First, the average \(^{90}\text{Sr}\) noise signal value was determined using multiple measurements of \(^{90}\text{Sr}\)-free NIST SRM-987. Then, the value is subtracted from \(^{90}\text{Sr}\) intensities of every measurement cycle in the run. All the data in the measurement were re-calculated off-line using the noise corrected \(^{90}\text{Sr}\) intensities, and finally, the noise corrected \(^{90}\text{Sr}/^{88}\text{Sr}\) ratios are determined. The correction value should be determined every analytical session because the noise intensity level is sensitive to the SEM parameters and analytical conditions. A single correction value is used for all measurements within the analytical session.

**\(^{90}\text{Sr}/^{88}\text{Sr}\) measurement of the reference materials and natural samples.** Table 2 summarizes the details of the analyzed environmental and biological samples, as well as the \(^{90}\text{Sr}\) activity parameters found in the literature (more details are shown in Table S1 in Supporting Information). As a \(^{90}\text{Sr}\)-free sample, a seawater reference material NASS-6 (Atlantic surface water) issued by the National Research Council Canada and a geochemical reference material JCp-1 (modern coral) issued by the Geological Survey of Japan were analyzed. The \(^{90}\text{Sr}\)-containing samples analyzed in this study include the International Atomic Energy Agency’s certified reference materials IAEA 156 (clover) and IAEA 330 (spinach). We also analyzed the ash from crayfish and smallmouth bass samples collected in Fukushima prefecture, Japan, after the Fukushima Daiichi Nuclear Power Plant accident. The \(^{90}\text{Sr}\) activity of these samples was analyzed using a radiometric method and published elsewhere.

Table 2 also summarized the measured Sr abundance of the \(^{90}\text{Sr}\) containing samples. The decay corrected \(^{90}\text{Sr}/^{88}\text{Sr}\) reference value of these samples was calculated for each analytical session using this measured Sr abundances together with either the certified \(^{90}\text{Sr}\) activity, for IAEA 156 and IAEA 330, or the radiometrically measured \(^{90}\text{Sr}\) activity, for Fukushima samples, respectively.

Table 2 \(^{90}\text{Sr}\) activity of the environmental and biological samples and results of Sr abundance measurement
| sample   | sample type     | $^{90}\text{Sr}$ [Bq/kg] | Sr [mg/kg]  |
|----------|-----------------|--------------------------|-------------|
| NASS-6   | seawater        | –                        | 7.48 $^{34}$ |
| JCp-1    | coral           | –                        | 7271 (218)  |
| IAEA 156 | clover          | 14.8 (3.0)               | 33.5 (2.0)  |
| IAEA 330 | spinach         | 20.1 (4.2)               | 44.0 (2.6)  |
| Ash 1    | crayfish        | 110 (8)                  | 1344 (66)   |
| Ash 2    | smallmouth bass | 25 (2)                   | 414 (28)    |

All the errors in the parenthesis are 2 SD.

The results of $^{90}\text{Sr}/^{88}\text{Sr}$ ratio measurements of the $^{90}\text{Sr}$ containing samples are summarized in Table 3 (more details are shown in Table S2 in Supporting Information), Figure 3, and Figure S5 in Supporting Information. The average $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of the $^{90}\text{Sr}$-free samples were -0.1 ± 2.5 × 10^{-12} (2SD, n = 5), 0.3 ± 2.9 × 10^{-12} (2SD, n = 8) and 0.5 ± 3.4 × 10^{-12} (2SD, n = 5) for NASS-6, JCp-1 (Fukushima-1) and JCp-1 (Fukushima-2), respectively. These results show that zero values of the $^{90}\text{Sr}/^{88}\text{Sr}$ ratios were measured without any systematic bias. No effect of sample-derived $^{90}\text{Zr}$ was observed down to 0.004 cps level. These results demonstrate the excellent stability of the noise signal, as well as the robustness of our noise correction scheme.

The measured $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of the IAEA 330 were 6.7 ± 1.3 × 10^{-11} (2SD, n = 4) and 5.76 ± 0.43 × 10^{-11} (2SD, n = 5) for temporally separated Bremen and Fukushima-2nd sessions. Both values agree well with the decay corrected reference $^{90}\text{Sr}/^{88}\text{Sr}$ ratios within analytical errors. The measured $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of the IAEA 156 were 3.26 ± 0.30 × 10^{-11} (2SD, n = 6) and it agrees well with the decay corrected reference value within analytical errors. The measured $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of the Ash1 (crayfish sample) was 1.38 ± 0.53 ×10^{-11} (2SD, n = 6) and 1.32 ± 0.42 ×10^{-11} (2SD, n = 6) for Fukushima-1st and Fukushima-2nd sessions, respectively. The measured $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of the Ash2 (smallmouth bass sample) were 1.40 ± 0.39 ×10^{-11} (2SD, n = 6) and 1.33 ± 0.27 ×10^{-11} (2SD, n = 6) for Fukushima-1st and Fukushima-2nd sessions, respectively. These measured values agree well with the decay corrected reference $^{90}\text{Sr}/^{88}\text{Sr}$ ratios within analytical errors (Table 3). No systematic bias exceeding the analytical uncertainties is observed for these results. All these measurements demonstrate the analytical capability of the energy-filtered TIMS technique, as well as our Sr separation chemistry to measure $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of 10^{-11} range (and possibly in the 10^{-12} range) in natural environmental and biological samples (Figure 4).

Table 3 Results of $^{90}\text{Sr}/^{88}\text{Sr}$ measurements of the environmental and biological samples
| sample      | reference\(^{90}\text{Sr}/^{88}\text{Sr}\) 1) | \(^{90}\text{Sr}/^{88}\text{Sr}\) 2) |
|------------|----------------------------------------|-------------------------------|
| NASS-6     | -                                      | -0.1 (2.5) × 10^{-12}         |
| JCp-1      | -                                      | 0.3 (2.9) × 10^{-12}          |
|            |                                        | 0.5 (3.4) × 10^{-12}          |
| IAEA 156   | 4.4 (0.9) × 10^{-11}                   | 3.26 (0.30) × 10^{-11}        |
| IAEA 330   | 7.9 (1.7) × 10^{-11}                   | 6.7 (1.3) × 10^{-11}          |
|            | 7.6 (1.7) × 10^{-11}                   | 5.76 (0.43) × 10^{-11}        |
| Ash 1      | 1.57 (0.15) × 10^{-11}                 | 1.38 (0.53) × 10^{-11}        |
|            | 1.56 (0.15) × 10^{-11}                 | 1.32 (0.42) × 10^{-11}        |
| Ash 2      | 1.16 (0.11) × 10^{-11}                 | 1.40 (0.39) × 10^{-11}        |
|            | 1.15 (0.11) × 10^{-11}                 | 1.33 (0.27) × 10^{-11}        |

1) Decay corrected to the reference date (see Table S2 in Supporting Information).

2) Noise corrected values using the average \(^{90}\text{Sr}\) count rate of NIST SRM 987 measurements (see text for details). All the errors in the parenthesis are 2 SD.

**Performance and applicability of \(^{90}\text{Sr}\) analysis by energy filtered TIMS.** Among the environmental and biological samples analyzed, Sr ion beams run short before 220 cycles in some of the measurements with a sample amount of 100 ng of Sr. Nevertheless, 100 ng measurements mostly show precision indistinguishable with the measurement of larger sample sizes. From these observations, we defined that the minimum amount of sample required for our \(^{90}\text{Sr}/^{88}\text{Sr}\) measurement as 100 ng of Sr. This is one of the essential parameters to evaluate the performance of TIMS measurement. Given the minimum sample size of 100 ng of Sr, and the \(^{90}\text{Sr}/^{88}\text{Sr}\) ratio detection limit of 2.7 ×10^{-12}, the detection limit of the absolute \(^{90}\text{Sr}\) amount can be estimated as 0.23 ag or 0.0012 mBq. The previous TIMS \(^{90}\text{Sr}\) measurements have reported a \(^{90}\text{Sr}\) detection limit of 0.17 fg (or 0.88 mBq)\(^{37}\) and 0.029 fg (or 0.15 mBq)\(^{38}\). Our technique has succeeded to lower the \(^{90}\text{Sr}\) detection limit of TIMS by three orders of magnitude.

Based on the applicability of this technique to environmental and biological sample measurements, detectable \(^{90}\text{Sr}\) activity concentrations in this study are different among samples with different stable Sr abundances (Figure 4). This is because the TIMS technique measures isotope ratios and thus the detection limit of the \(^{90}\text{Sr}\) analysis is primarily determined as \(^{90}\text{Sr}/^{88}\text{Sr}\) ratio. For example, seawater is
one of the typical environmental samples and has a stable Sr abundance of approximately 8 mg/kg. With this energy-filtered TIMS technique, the detection limit of seawater $^{90}$Sr activity concentration is estimated as 0.09 Bq/kg. The minimum quantity of the seawater sample required for the measurement is as small as 12.3 μL. River water and groundwater have stable Sr abundances range from several tens to several hundreds of ppb. In this case, the detection limit of $^{90}$Sr activity concentration will be as low as 0.1 mBq/kg. A sample size of 10 g will be required to analyze such a dilute sample.

Comparing the capability of $^{90}$Sr analysis between TIMS and ICP-MS is not straightforward. The capability of ICP-MS is limited by the detection limit of the absolute amount of $^{90}$Sr and the abundance sensitivity defined as the $^{90}$Sr/$^{88}$Sr ratio (Figure 5). It is not limited by the amount of stable Sr unless the $^{90}$Sr/$^{88}$Sr ratio was interfered with the abundance sensitivity limit. However, the capability of TIMS is limited by the detection limit of $^{90}$Sr/$^{88}$Sr ratio and the minimum amount of sample required for the analysis (Figure 5). The latter is required to keep the fixed ion beam intensity of $^{88}$Sr during the measurement because the performance of the $^{90}$Sr/$^{88}$Sr ratio measurement by TIMS is mainly determined by the ion beam intensity of the most abundant $^{88}$Sr. The detection limit of the absolute amount of $^{90}$Sr is a derivative of these two parameters for TIMS. Compared with the previous TIMS and ICP-MS/MS techniques, a significant reduction in the $^{90}$Sr/$^{88}$Sr detection limit achieved in this study had widely expanded the applicability of this technique to analyze samples with low $^{90}$Sr contents (Figure 5). Additionally, an order of magnitude reduction of the minimum sample amount compared with the previous TIMS technique allows the analysis of 10 times smaller sample sizes. These features are advantageous in analyzing trace amounts of $^{90}$Sr on size-limited environmental samples such as soil exchangeable fraction and eolian dust as well as biological samples with low Sr abundances.

**Methods**

**Reagents and standards.** For sample preparation, ultrahigh purity grade acids, HNO$_3$, HCl, HClO$_4$, and HF (TAMAURE AA-100; Tama Chemicals Co., Ltd, Japan.), as well as ultra-pure water produced by a Milli-Q Element system (Millipore, USA) were used. The Sr isotopic reference material SRM 987 issued by NIST was used as a $^{90}$Sr-free reference standard.

**Sample preparation.** The IAEA reference material’s certified $^{90}$Sr activity values were calculated using their dry weight. Therefore, these materials were weighed and digested as they were. Concentrated HNO$_3$ and HClO$_4$ acids were used to digest IAEA 156 and IAEA 330. The Fukushima samples were analyzed on an ashed sample, which were used in the previous $^{90}$Sr activity analysis. The ashed crayfish and smallmouth bass samples were weighed and digested successively using 2.4M HCl and conc HNO$_3$. JCp-1 was digested with 5% CH$_3$COOH. Finally, all the samples were dissolved in 3M HNO$_3$. NASS-6 was mixed with 6 M HNO$_3$ to make a 3 M HNO$_3$ solution.
Sr was separated from the other elements, including Zr, using extraction chromatography with Sr Resin (Eichrom Technologies Inc., USA) using a handmade PTFE column of 0.2 mL volume. After sample solution loading, the column was rinsed with 2 mL of 6 M HNO₃ and 0.5 mL of 3 M HNO₃ successively, and Sr was eluted with 2 mL of 0.05 M HNO₃. Finally, the separated Sr fraction was reacted with one drop of concentrated HNO₃ to decompose resin-derived organics. The total yield of the Sr separation chemistry was approximately 94%. The elution curve estimated the decontamination factor of Zr against Sr in our column chemistry at 1.2 × 10⁻⁵.

Quantitative analysis of Sr and the reference $^{90}$Sr/$^{88}$Sr ratio of the samples. The Sr concentrations of the $^{90}$Sr containing samples (IAEA 156, IAEA 330, crayfish, and smallmouth bass) were measured using ICP-MS Agilent 7700 (Agilent Technologies, Japan) and a sample of a fraction of the sample solution at Kochi Institute for Core Sample Research, JAMSTEC. The measured Sr concentration was used to calculate the reference $^{90}$Sr/$^{88}$Sr values.

Instrumentation. In this study, two thermal ionization mass spectrometers were used, the Triton™ XT (Thermo Scientific™) at Thermo Fisher Scientific, Bremen, Germany, and the Triton™ Plus (Thermo Scientific™) at Fukushima University, Japan. Both instruments are equipped with eight movable Faraday cups, an axial Faraday cup, and an axial secondary electron multiplier (SEM). The Faraday cups are connected with the standard $10^{11}$ Ω amplifiers. A set of RPQ lenses, which is an energy filtering device, is equipped in front of the SEM detector. The configuration of the detectors was summarized in Table S3 in supporting information.

This study uses both Re and Re-ZR (rhenium zone-refined) single filaments. High temperature bakeout of the filaments at 5.5 A filament current were performed under high vacuum for 30 min (Bremen) or 60 min (Fukushima). In Fukushima, the baked filaments are further inspected for $^{90}$Zr emission. $^{90}$Zr intensity of the baked filaments is measured at filament temperature of approximately 1550°C for 1200 s. Filaments with $^{90}$Zr intensities higher than 0.04 cps are disposed of at this point.

$^{90}$Sr/$^{88}$Sr ratio measurement by TIMS.

The separated sample containing 100–400 ng of Sr was loaded onto a baked filament with a Ta activator. The amount of Sr in the samples was fixed just before the Sr separation. The actual sample amount on the filament is lower than this nominal amount because losses of sample during the Sr separation and sample loading are not included in this number. In Fukushima, dull-red heating during sample loading was avoided because the sample becomes fragile and tends to be partially lost from the
filament surface. In both laboratories, Sr isotopes are measured using static multcollection mode with the slightly shifted axial mass settings of 89.8777 to eliminate the isobaric interference of organic molecules on m/z 90.0. In Bremen, filaments were initially heated to 2400 mA at a rate of 300 mA/min while measuring a 5-min baseline for Faraday cup detectors. The filament current increased at 100 mA/min until an $^{88}\text{Sr}$ signal of 20 V was achieved. After a 1-min wait, the filament current was re-adjusted and the measurement was initiated. Ion beam integration of 4.194 s was repeated for approximately 1 h (10 blocks, 86 cycles per block). During the measurement, the filament current was adjusted to keep the $^{88}\text{Sr}$ signal within 20 ± 4 V. In Fukushima, filaments were preheated at 2250 mA for 30 min before the measurement. During the measurement, filaments were heated to 1500 mA at a rate of 1500 A/min and then to 2250 mA at a rate of 214 mA/min. The filament current was increased at 100 mA/min until an $^{88}\text{Sr}$ signal of 25 V was achieved. Ion beam integration of 16.777 s was repeated for approximately 1 h (22 blocks, 10 cycles per block). During the measurement, the inter-block filament current adjustment kept the $^{88}\text{Sr}$ signal at ± 2.5 V. The Faraday cup baseline was measured using a long baseline method utilizing the very stable nature of the Faraday cup noise level of TRITON. A 1 h baseline was measured for every 10 h of the measurement sequence. After the first analytical session in Fukushima, the operating voltage of the SEM detector was adjusted, resulting in slightly different analytical conditions between the two sessions. In both laboratories, FC-SEM yield was calibrated just before the measurement. The dark noise of the SEM detector was measured for 1h along with the Faraday cup baselines. However, all the measurements are not corrected for dark noise. Instead, the average value of multiple $^{90}\text{Sr}$-free NIST 987 measurements was used to correct overall noise signals on $^{90}\text{Sr}$. The analytical error of the $^{90}\text{Sr}/^{88}\text{Sr}$ ratio was determined as the 2σ internal error, which is two times the standard error of the $^{90}\text{Sr}/^{88}\text{Sr}$ ratios of 220 or 860 cycles. In all the measurements, the analytical error of the $^{90}\text{Sr}/^{88}\text{Sr}$ ratio was comparable with the $\sqrt{n}$ counting error of the accumulated $^{90}\text{Sr}$ count before noise correction.

Instrumental isotope fractionation is a possible error source of the $^{90}\text{Sr}/^{88}\text{Sr}$ ratio measured using mass spectrometry. However, the variability of the average $^{88}\text{Sr}/^{86}\text{Sr}$ values caused by the instrumental mass fractionation during the measurement was ± 0.5% ($^{88}\text{Sr}/^{86}\text{Sr}$ values from 8.31 to 8.41) and was negligible compared with the analytical error of $^{90}\text{Sr}/^{88}\text{Sr}$ ratio (> 10 %, 2RSD) in this study. Therefore, no correction for instrumental isotope fractionation was made.

**Conclusions**

This study demonstrates the ability of a TIMS instrument to perform an accurate $^{90}\text{Sr}/^{88}\text{Sr}$ ratio measurement of samples with $^{90}\text{Sr}/^{88}\text{Sr}$ ratios as low as $2.7 \times 10^{-12}$. Not only the introduction of effective hardware such as the RPQ lenses but a thorough investigation and elimination of noise signals have resulted in a constantly low noise signal of $^{90}\text{Sr}$ on the 0.001 cps order. The sample throughput of this technique was 40 samples per day for Sr separation chemistry and 21 samples per 2 days for isotope ratio measurement. This is not as fast as the ICP-MS techniques but is significantly faster than radiometric methods. To detect 0.0012 mBq of $^{90}\text{Sr}$, this energy-filtered TIMS technique requires a sample
size of 100 ng of Sr. This sample size corresponds to 12.3 μL of seawater or less than 10 mg of biological samples with Sr abundances larger than 10 mg/kg, which seems easy to handle. The lowered $^{90}$Sr detection limit and smaller sample size of this technique are suitable for studying environmental diffusion of radioactive materials, as well as environmental elemental cycling studies using nuclear test origin $^{90}$Sr in the earth surface materials.

**Declarations**

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**Supporting Information**

The Supporting Information is available free on the *Scientific Reports* website at http://*****.

Detailed experimental section: $^{90}$Sr activity parameters of the environmental samples (Table S1); Transition of Noise signal during Sr isotope ratio measurement (Figure S1); Measured $^{90}$Sr/$^{88}$Sr ratios of NIST SRM987 (Figure S2); Measured $^{90}$Sr/$^{88}$Sr ratios of the environmental samples (Figure S3); $^{90}$Sr/$^{88}$Sr ratios of the environmental samples (Table S2) (PDF).

**Author Contributions**

S.W.: Method set-up, investigation and writing paper draft. J.A. and R.S.: Sample measurement in Fukushima, investigation. K.S: discussion, writing —review and editing—. T.M: discussion, writing —review and editing—. J.R: Method set-up, Sample measurement in Bremen, data evaluation, discussion. H.V: Method set-up, data evaluation and discussion. S.S: discussion. Y.T.: conceptualization, supervision, and project administration.

**Notes**

The authors declare no competing financial interest.

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**Figures**

![Graph showing noise signal intensity on ⁹⁰Sr [cps] versus ⁸⁸Sr [V] with different conditions: without RPQ, RPQ installed, and RPQ fine tuned.](image-url)
Figure 1

Noise signal intensity on 90Sr was plotted against the intensity of 88Sr. All data were taken using a sample of 90Sr-free NIST SRM987. Black triangle represents the data taken without RPQ. The blue diamond represents the data taken with RPQ but without fine-tuning. Red circle represents the data taken with fine-tuned RPQ. The positive correlation of the data shows the effect of the 88Sr peak tail on 90Sr.

Figure 2
Results of repeated 90Sr measurements of 90Sr-free NIST SRM987 in (a) Bremen, (b) Fukushima-1, and (c) Fukushima-2 sessions. White diamonds represent noise uncorrected data of a single measurement. Blue line and blue bands represent the average and 2 SD range of the multiple measurements. The grey area represents the average intensity of dark noise observed during each measurement session.

The measured 90Sr/88Sr ratio of the environmental and biological samples was compared with the decay corrected reference 90Sr/88Sr ratios. The blue circle represents an average of multiple measurements of the sample. Error bars correspond to 2 SD.
Stable Sr and 90Sr abundances of the analyzed samples (blue circles) were compared with the 90Sr/88Sr detection limit (red line) of this study. The 90Sr/88Sr detection limit of the previous TIMS studies was plotted for comparison. Stable Sr abundance range of selected environmental samples is also plotted: plants, seawater, and river water/groundwater.
Figure 5

The capability of the mass spectrometric 90Sr determination methods is demonstrated using the absolute amount of stable Sr and 90Sr required for a single measurement. Samples within the shaded area are measured by the method. The lowest point of the shaded area corresponds to the detection limit of 90Sr for TIMS.

Supplementary Files

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