Cell Biometrics Based on Bio-Impedance Measurements

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1. Introduction

Many biological parameters and processes can be sensed and monitored using their impedance as marker (Grimmes, 2008), (Beach. 2005), (Yúfera, 2005), (Radke, 2004), with the advantage that it is a non-invasive, relatively cheap technique. Cell growth, cell activity, changes in cell composition, shapes or cell location are only some examples of processes which can be detected by microelectrode-cell impedance sensors (Huang, 2004) (Borkholder, 1998). The electrical impedance of a biological sample reflects actual physical properties of the tissue. In frequency dependent analyses, the $\beta$-dispersion ranging from kilohertzs to hundreds of megahertzs (Schwan, 1957) is mainly affected by the shape of the cells, the structure of the cell membranes, and the amount of intra and extra cellular solution. Electrical bio-impedance can be used to assess the properties of biological materials (Ackmann, 1993) involved in processes such as cancer development (Giaever, 1991), (Blady, 1996), (Aberg, 2004); because the cells of healthy tissues and cancer are different in shape, size and orientation, abnormal cells can be detected using their impedance as a marker.

Among Impedance Spectroscopy (IS) techniques, Electrical Cell-substrate Impedance Spectroscopy (ECIS) (Giaever, 1986), based on two-electrode setups, allows the measurement of cell-culture impedances and makes it possible to determine the biological condition (material, internal activity, motility and size) of a cell type and its relationship with the environment; for example, the transfer flow through the cell membrane (Wang, 2010). One of the main drawbacks of the ECIS technique is the need to use efficient models to decode the electrical results obtained. To efficiently manage bio-impedance data, reliable electrical models of the full system comprising electrodes, medium and cells are required. Several studies have been carried out in this field (Giaever, 1991), (Huang, 2004), (Borkholder, 1998), (Joye, 2008), (Olmo, 2010), some of them employing Finite Element simulation (FEM) for impedance model extraction. These models are the key for matching electrical simulations to real system performances and hence for correctly decoding the results obtained in experiments.
The use of FEM analysis with programs such as FEMLAB (Huang, 2004) considers that the DC mode can be employed for sinusoidal steady-state calculation by assigning complex conductivity. This works because the Poisson equation has the same form as the Laplace equation in the charge-free domain. In this chapter we will describe an alternative method of performing FEM simulations of electrode – cell interfaces based on COMSOL. The quasistatic mode of COMSOL is used, which also takes into account magnetic fields to calculate electrical impedance. The models obtained are successfully applied as loads for full electrical system modelling and simulation. In this sense, any biological sample could be electrically modelled and simulated, facilitating reliable information about integrated circuit design for impedance measuring (Yúfera, 2010b). This work includes several improvements to the model in (Huang, 2004), both to the cellular membrane and to the cell-electrode gap region. Impedance changes on small electrodes (32 µm square) caused by 3T3 mouse fibroblasts are simulated in order to validate the model and characterize the microelectrode sensor response to cell size and growth.

The knowledge acquired from the electrode-cell model can be used to create a set of applications useful in cell culture biometry and for improving efficiency biology lab tasks. One use is the detection of cell sizes by characterizing the model in terms of cell-area overlap. Impedance sensor sensitivity curves with cell size will be presented. As an extension of the application for the aforementioned sensitivity curves, a technique for measuring the cell index (CI) coefficient in cell-culture growth processes will be presented and we will show how, as a consequence of this, cell toxicity experiments can be monitored in real-time. An Analog Hardware Description Language (AHDL) model (SpectreHDL) for the mixed-mode simulation of full (electronic and biological) systems will also be described. By applying the developed model to ECIS curves obtained experimentally, it will be possible to determine cell density and toxin-caused cell death rates. Moreover, cell modelling has recently been applied to cell imaging or bioimpedance microscopy. In (Linderholm, 2006) Electrical Impedance Tomography (EIT) techniques were reported for cell motility detection. The models proposed can also be applied to decode impedance measurements obtained from cell culture measurements, producing a two dimensional cell location map; that is, a microscopy image based on bio-impedance measurements (Yúfera, 2011).

Having explained the objectives of the work, the proposed contents of this chapter are as follows. The second section gives a brief overview of electrode solution models useful for cell-electrode characterization. The third section presents a useful method for generating cell-electrode electrical models based on COMSOL multiphysics software. The fourth section describes the finite element simulations performed. Processes for extracting useful models are included in the fifth section, which also illustrates cell size detection simulations on a simplified system. AHDL models are presented in sixth section. The seventh section covers the real time monitoring of the cell under cultivation and the application of the proposed model in dosimetric experiments. Finally, in the eighth section, a two dimensional approach to bioimpedance microscopy is described, based on the models previously developed. Conclusions are given in the ninth section.

2. The electrode – electrolyte electrical model

The impedance of electrodes in ionic liquids has been researched quite extensively (Robinson, 1968), (Schwan, 1963), (Simpson, 1963), (Schwan, 1992), (Onaral, 1982) and (Onaral, 1983). When a solid (including metals, semiconductors, and insulators) is immersed
in an ionic solution, ions in the solution can react with the electrode and the solid ions from the electrode can enter the solution, leading to complex reactions at the interface. An electrified interface or double layer develops at the interface of the two phases. Eventually, electrochemical equilibrium is established at the interface: the current flowing into the electrode is equal but opposite in sign to that flowing out from the electrode. The net result is the development of a charge distribution at the interface with an associated electric potential distribution. The Helmholtz-Gouy-Chapman-Stern model is the commonly accepted model for describing the charge distribution at the electrode interface (Bockris, 1970).

The negatively charged electrode attracts hydrated ions with positive charges to the surface but repels negatively charged ions away from the surface, yielding the profiles of cation and anion concentration $C^+$ and $C^-$, respectively. The water dipoles are also reoriented under the electric fields. Some ionic species that are not obstructed by their primary hydration sheath, such as some anions, can make their way to and come into contact with the electrode. Most cations have a water sheath due to their lower dissolvation energy (Bockris, 1970) and a smaller contact angle with water dipoles. The charge distribution extends to the bulk solution thanks to thermal motion, forming an ion cloud-like diffusion layer and a charge spatial distribution. The profile of the diffuse zone depends on the Debye length, which in turn depends on the gas constant, the temperature, the ion charge number and the ion concentration of the bulk solution. (Borkholder, 1998).

When a sufficiently small sinusoidal current is applied to the electrode at equilibrium, the electrode potential will be modulated by a sinusoidal overpotential (Schwan, 1968). In the range of linear behavior, the phasor ratio of the output overpotential to the input current defines the AC polarization impedance. During the small current perturbation, charge transfer due to chemical reactions and mass diffusion all occur at the electrode surface. The rate determining step will dictate the electrode polarization impedance.

In the following paragraph, we will first discuss an equivalent circuit representing all the phenomena occurring at the electrode-solution interface and then we will explain each component in the circuit. The electrified interface can be considered as the series connection of two parallel-plate capacitors with the thicknesses of a compact layer and a diffuse layer respectively and with a water dielectric. This is the electrode-solution interfacial capacitance of the electrified double layer. Apart from the double layer capacitance $C_I$, the electrode-solution interface has faradic impedance representing a barrier to current flow between the electrode and the solution, including the finite rate of mass transport and electron transfer at the electrode surface. These phenomena are modelled in the equivalent circuit in Fig. 1, in which the faradic impedance is in parallel with the double layer capacitance.

The current flowing through the electrified interface will encounter a resistance $R_{ct}$ caused by the electron transfer at the electrode surface and Warburg impedance $Z_W$ due to limited mass diffusion from the electrode surface to the solution. As a result, in the equivalent, the electron transfer resistance $R_{ct}$ is in series with the mass diffusion limited impedance $Z_W$. As the current spreads to the bulk solution, the electrode has a solution conductivity-determined series resistance, represented as spreading resistance $R_S$ in the equivalent circuit.

### 2.1 Double layer capacitance ($C_I$)

The region between the electrode surface and the Outer Helmholtz Plane (OHP) consists mostly of water molecules (Borkholder, 1998). The thickness of the OHP layer is $x_{OH}$ (distance from the metal electrode to OHP), and consequently the capacitance of the Helmholtz layer $C_{II}$ is given by,
Fig. 1. Equivalent circuit for electrode solution interface. $C_I$ is the double layer capacitance, Faradic impedance includes $Z_w$, the Warburg impedance and $R_{ct}$, the charge-transfer resistance. $R_s$ is the spreading resistance.

$$C_H = \varepsilon \varepsilon_0 / \varepsilon_H$$  \hspace{1cm} (1)

The capacitance of the diffuse layer can be derived by differentiating the surface charge with respect to the potential. The total interfacial capacitance $C_I$ of the electrified double layer consists of the series combination of the Helmholtz compact layer and the diffuse layer.

$$1 / C_I = 1 / C_H + 1 / C_G$$  \hspace{1cm} (2)

where $C_H$ is the capacitance of the Helmholtz layer described earlier and $C_G$ is the Gouy-Chapman capacitance due to the diffuse ion cloud. This double layer capacitance has been studied in many works (Borkholder, 1998), (Schwan, 1968), (Simpson, 1980), (De Boer, 1978) although most of them agree on the value of $C_I$ (approximately 15 µF/cm$^2$).

### 2.2 Warburg impedance ($Z_w$)

Warburg impedance is related to the mass diffusion process occurring in the electrode-electrolyte interface (Warburg, 1899). In AC measurements, in response to the sinusoidally varying potential, the ion concentration gradient at the interface increases with frequency and the ions diffuse less as frequency increases. The Warburg impedance follows this expression,

$$Z_w = \omega^{-1/2} K_w / A_e (1 + j)$$  \hspace{1cm} (3)

where $A_e$ is electrode area and $K_w$ [Ω sec$^{-1/2}$·cm$^2$] is a constant determined by the electrochemistry and mobility of the ions involved in the charge transfer reaction. It is difficult to find a theoretical value for $K_w$, which in some works seems to be a parameter that is included in the model to adapt the model to experimental measurements (Huang, 2004). In other works, this impedance is not taken into account, because it is considered negligible for the materials and frequency range used in electrophysiological experiments (Joye, 2008).

D. A. Borkholder, in his thesis, gives some reference values for $Z_w$ for circular bare platinum electrodes of different sizes (Borkholder, 1998).

### 2.3 Charge transfer resistance ($R_{ct}$)

Charge transfer resistance is determined by the electron transfer rate at the interface. Electrodes made of noble metal electrodes such as platinum, gold etc., in physiological saline solution act as catalytic surfaces for the oxygen redox reactions (De Rosa, 1977), (Bagotzky, 1970). The reaction rates for the anodic and cathodic processes are not the same. In the state of equilibrium, the interface current due to oxidation is equal but of the opposite
Cell Biometrics Based on Bio-Impedance Measurements

347

sign to that caused by reduction. Additional potential applied to the electrode will cause a net current flow. In the range comparable to the equilibrium current, the electrode behaves as a linear resistive component the resistance of which is referred to as the charge transfer resistance. As the voltage increases, the excess current increases exponentially and the charge transfer resistance decreases exponentially with the applied voltage. The charge transfer resistance can be described as,

\[ R_{ct} = \frac{V_t}{J_0 z} \]  

(4)

where \( J_0 \) is the exchange current density (A/cm\(^2\)), \( V_t \) is the thermal voltage (K\(T/q\)) and \( z \) is the valence of the ion involved in the charge transfer reaction.

2.4 Spreading resistance (\( R_s \))

The final circuit element which must be included in the basic electrode/electrolyte model is the spreading resistance. As the name implies, this resistance models the effects of the spreading of current from the localized electrode to a distant counter electrode in the solution. It can be calculated by integrating the series resistance of solution shells moving outward from the electrode, where the solution resistance (\( R \) in \( \Omega \)) is determined by

\[ R_s = \rho L/A \]  

(5)

where \( \rho \) is the resistivity of the electrolyte (\( \Omega \* cm \)), \( L \) is the length between sensing and counting electrodes (cm) and \( A \) is the cross-sectional area (cm\(^2\)) of the solution through which the current passes. A similar expression is used in most models (Borkholder, 1998), (Joye, 2008).

3. Finite Element Model (FEM)

3.1 Cell electrode model

We first explored the work performed by Huang et al., making use of the computational advantages offered by COMSOL (http://www.comsol.com) over FEMLAB. Our objective was to obtain a model for the impedance changes caused by cell growth on electrodes similar in size to the cell. The cells modeled in the simulation were 3T3 mouse fibroblasts, which closely attach to surfaces and which typically have a cell-surface separation 0.15 \( \mu \)m. The cells are about 5 \( \mu \)m in height and, seen from above, are irregularly shaped and approximately 30–50 \( \mu \)m in extent. A circular cell 30 \( \mu \)m in diameter centered on a 32x32 \( \mu \)m square sensing electrode was considered. (see Fig. 2). The sensing electrode was surrounded by a counter electrode with a considerably greater area.

3T3 mouse fibroblasts consist of a thin (about 8 nm), poorly conducting membrane surrounding a highly conductive interior (Giebel, 1999). The cell culture medium simulated by Huang et al. is highly ionic and possesses a conductivity of approximately 1.5 S/m. The cell culture medium fills the cell-electrode gap and forms an electrical double layer (Helmholtz plus diffuse layer) approximately 2 nm thick between the bulk of the medium and the electrode.

Some approximations were made by Huang et al. to address the problem using FEMLAB. Only one quarter of the electrode was simulated. As the problem involves a wide range of distance scales, it was difficult to solve by finite-element techniques, so the following adjustments were made:
The electrical double layer modelling the electrode-solution equivalent circuit was replaced with a 0.5 µm thick region with the same specific contact impedance.

\[
\sigma_{dl} + j 2\pi f \varepsilon_{dl} = t \left[ \frac{(2\pi f)^{0.5}}{K_w} + j \frac{(2\pi f)^{0.5}}{K_w} + j C_I 2\pi f \right]
\]  

(6)

where \(\sigma_{dl}\) and \(\varepsilon_{dl}\) are the conductivity and dielectric permittivity of the double layer, \(t\) is the thickness of the region, \(C_I\) is the interfacial capacitance per unit area, comprising the series combination of the Helmholtz double layer and the diffuse layer, and \(K_w\) is a constant related to the Warburg impedance contribution.

The cell membrane was replaced by a 0.5 µm thick region with the same capacitance per unit area,

\[
c_m = t C_m
\]

(7)

where \(C_m\) is the membrane capacitance per unit area and \(t = 0.5\) µm.

The electrode-cell gap was replaced with a 0.5 µm thick region with the same sheet conductivity, that is

\[
\sigma_{gap} = \frac{t_{cell-electrode}}{t} \sigma_{medium}
\]

(8)

where \(t_{cell-electrode}\) is the gap thickness and \(t\) is again 0.5 µm.

In our study we adopted the geometry of their simulation (see Fig. 2), and the values for the conductivity and permittivity of the electrical double layer were calculated following the same expression as shown eq. (6), with the same values for \(K_w\) and \(C_I\) as those mentioned in (Huang, 2004). Conductivity of the cell and the medium was also set to 1.5 S/m in our work. However, the model used by X. Huang et al. for the electrode-cell gap and the cellular membrane (eqs. 7 and 8) was refined as shown in the following section.

### 3.2 Model enhancement

Several modifications were made in the model in order to simulate cell impedance measurements more accurately and obtain a more complex model that reflected real experiments in a more realistic way. These modifications were made in the following areas.
3.2.1 Cell membrane

The equivalent circuit of the attached membrane was modeled as a resistance $R_m$ in parallel with a capacitance $C_m$, in a similar way to that reported in (Borkholder, 1998). These parameters are defined as,

$$
\frac{1}{R_m} = \frac{1}{g_{\text{mem}} A_e}
$$

$$
C_m = \varepsilon_m A_e
$$

where $A_e$ is the area of the attached membrane (in our case $A = 706.86 \times 10^{-12} \text{m}^2$), $g_{\text{mem}} = 0.3 \text{mS/cm}^2$ is the local membrane conductivity and $c_{\text{mem}} (1 \mu\text{F/cm}^2)$ is the membrane capacity per unit area. We can calculate the conductivity and permittivity of the cellular membrane from the impedance using the following expression

$$
Z = \frac{1}{K (\sigma + j\omega \varepsilon)}
$$

where $K$ is the geometrical factor ($K = \text{area} / \text{length}$). In our case a value of 5 $\mu$m was taken as the length. (This value corresponds to the thickness of the membrane layer represented in COMSOL). The value obtained for $K$ was $1.413 \times 10^{-3}$, and the values obtained for conductivity and permittivity were $\sigma = 1.5 \mu\text{S/m}$ and $\varepsilon = 5.001 \text{nF/m}$ ($\varepsilon_r = 565$).

3.2.2 Cell membrane-electrolyte interface capacitance

This capacitance was not considered in Huang’s model, but may also be important, as it models the charge region (also called the electrical double layer) which is created in the electrolyte at the interface with the cell. The capacitance $C_{hd}$ is defined as the series of three capacitances,

$$
C_{h1} = \frac{\varepsilon_0 \varepsilon_{\text{IHP}}}{d_{\text{IHP}}} A_e
$$

$$
C_{h2} = \frac{\varepsilon_0 \varepsilon_{\text{OHP}}}{d_{\text{OHP}} - d_{\text{IHP}}} A_e
$$

$$
C_d = \frac{q \sqrt{2 \varepsilon_0 \varepsilon_r K T z n_0 N}}{K T} A_e
$$

where $A_e$ is the area of the attached membrane, $\varepsilon_0$ is the dielectric permittivity of free space; $\varepsilon_{\text{IHP}}$ and $\varepsilon_{\text{OHP}}$ are, respectively, the Inner and Outer Helmholtz Plane relative dielectric constant; $d_{\text{IHP}}$ is the distance from the Inner Helmholtz Plane to the membrane; $d_{\text{OHP}}$ is the distance from the Outer Helmholtz Plane to the membrane; $\varepsilon_d$ is the diffuse layer relative dielectric constant; $K_B$ is Boltzmann’s constant; $T$ is the absolute temperature; $q$ is the electron charge; $z$ is the valence of ions in the solution; $n_0$ is the bulk concentration of ions in the solution; and $N$ is the Avogadro constant.

For $C_{hd}$, the values given in (Joye, 2008) are considered. In particular, it is assumed that $\varepsilon_{\text{IHP}} = 6$, $\varepsilon_{\text{OHP}} = 32$, $d_{\text{IHP}} = 0.3 \text{nm}$, $d_{\text{OHP}} = 0.7 \text{nm}$, $z = 1$, $T = 300 \text{K}$, and $n_0 = 150 \text{mM}$. The area of the attached membrane is in our case $A_e = 706.86 \mu\text{m}^2$ and $\varepsilon_d$ is set to 1. The following values were obtained: $C_{h1} = 0.125 \text{pF}$; $C_{h2} = 0.5 \text{pF}$; $C_d = 2.22 \text{pF}$, and the total series capacitance was
$C_{hd} = 0.1 \text{pF}$. Comparing the impedance equivalent to this capacitance with the same expression as before (eq. 10), and remodelling this layer as a 5 µm thick layer with $K = 1413.10^{-6}$, we obtained $\varepsilon = 0.0011 \text{uF/m}$, which corresponds to $\varepsilon_r = 124.29$, the value that was loaded into COMSOL.

4. COMSOL Simulations

As can be seen in Fig. 2, only one quarter of the electrodes and the cell was simulated. Electrodes were modeled with no thickness. The first layer modeled on top of the electrode surface was the 0.5 µm thick electrical double layer, which can be seen in the figure. On top of the electrical double layer, the cell-electrode gap was modeled with another 0.5 µm layer. In our simulation this layer included the cell membrane-electrolyte interface capacitance. Finally, on top, we have the cell membrane, also modeled as another 0.5 µm layer, and the rest of the cell. For each layer, it is necessary to load the conductivity and permittivity values calculated earlier into COMSOL. All surfaces had an insulating boundary condition ($n*J=0$) with the exception of the surfaces separating the different layers and sub-domains within the model, which were set to continuity ($n.(J_1-J_2) = 0$), and the bottom surfaces of the two electrodes, which were set to an electric potential of 1V and 0V.

The Quasi-statics module of COMSOL was used to perform the finite element simulations. In this mode, it is possible to obtain the solution for the electric potential for different frequencies. The simulations were performed on a 2.26 GHz Intel(R) Core(TM)2 DUO CPU. Solution times varied with the frequency but ranged from 3 to 6 minutes. In Fig. 3 we can see the solution for the electric potential at the determined frequency of 100 Hz.

Fig. 3. Electric potential solution at 100Hz.

Two series of simulations, at frequencies ranging from $10^2$ Hz to $10^6$ Hz, were made with and without the presence of the cell. Once the solution for the electric potential had been found by COMSOL, Boundary Integration was used to find the electric current through the counter electrode. With that value the electric impedance was calculated, taking into account that the voltage difference between electrodes was 1V and that impedance had to be divided by 4 (as only one quarter of the electrodes was being simulated.) The values obtained are shown in Fig. 4.
The measured impedance changes by several orders of magnitude over the simulated frequency range, in accordance with previous works (Schwan 1992), (Onaral, 1982). It can be appreciated how the presence of the cell changes the measured impedance, with the biggest change recorded at a frequency near to $10^5$ Hz. This also corroborates the study carried out by (Huang, 2004). Another way of representing the impedance magnitude is to observe the impedance changes in the system with the cell on top with respect to the microelectrode system without cell. This can be done by plotting the normalized impedance change with respect to the cell-less system, defined as

$$r = \frac{Z_c - Z_{nc}}{Z_{nc}}$$

(12)

where $Z_c$ and $Z_{nc}$ are the impedance magnitudes with and without cell, respectively. The normalized impedance changes in the system with the 30 µm-diameter cell modelled earlier is plotted in Fig. 4 (red line). Figure 5 shows the normalized impedance parameter, $r$, as a function of frequency. It can be seen how the 100 kHz frequency seems to be the optimal value for sensing the cell, since sensor sensitivity is maximum.

Fig. 4. Impedance magnitude of the microelectrode system with (red line) and without (blue line) the cell.

Fig. 5. Simulated normalised impedances of the system, for a 30 µm cell diameter.
5. Model extraction process

Figure 6 shows an example of a two-electrode impedance sensor useful for the ECIS technique: \( e_1 \) is called the sensing electrode and \( e_2 \) the reference electrode. Electrodes can be manufactured in CMOS process with metal layers (Huang, 2004) or using post-processing steps (Manickam, 2010). The cell location on top of \( e_1 \) must be detected. Circuit models developed to describe electrode-cell interfaces (Huang, 2004, Joye, 2008) contain technology process information and take the cell-electrode overlap area as their main parameter. The correct interpretation of these models gives information for:

a) **electrical simulation**: parametrized models can be used to update the electrode circuit in terms of its overlap with cells,

b) **imaging reconstruction**: electrical signals measured with sensors can be associated with a given overlap area, making it possible to obtain the actual covering on the electrode from experimental results.

![Fig. 6. Basic concept for measuring with the ECIS technique using two electrodes: \( e_1 \) (sensing electrode) and \( e_2 \) (reference electrode). AC current \( i_x \) is applied between \( e_1 \) and \( e_2 \), and voltage response \( V_x \) is measured from \( e_1 \) to \( e_2 \), including the effect of \( e_1 \), \( e_2 \) and sample impedances.](image)

This work employs the electrode-cell model reported in (Huang, 2004, Olmo, 2010), obtained using finite element method simulations. The model in Fig. 7 considers that the sensing surface of \( e_1 \) could be total or partially filled by cells. For the two-electrode sensor in Fig. 6, with \( e_1 \) sensing area \( A \), \( Z(\omega) \) is the impedance by unit area of the empty electrode (without cells on top). When \( e_1 \) is partially covered by cells in a surface \( A_c \), \( Z(\omega)/(A-A_c) \) is the electrode impedance associated with the area not covered by cells, and \( Z(\omega)/A_c \) is the impedance of the area covered. \( R_{gap} \) models the current flowing laterally in the electrode-cell interface, which depends on the electrode-cell distance at the interface (in the range of 15-150 nm). \( R_s \) is the spreading resistance through the conductive solution. For an empty electrode, the impedance model \( Z(\omega) \) is represented by the circuit in Fig. 1. For \( e_2 \), not covered by cells, the model in Fig 7a was considered. The \( e_2 \) electrode is typically large and grounded, and its resistance is small enough to be rejected. Figure 8 represents the impedance magnitude, \( Z_c \), for the sensor system in Fig. 6, considering that \( e_1 \) could be either empty, partially or totally covered by cells. The parameter \( ff \), called the **fill factor**, can be zero for \( A_c = 0 \) (\( e_1 \) electrode empty), and 1 for \( A_c = A \) (\( e_1 \) electrode full). \( Z_c (ff=0)= Z_{nc} \) is defined as the impedance magnitude of the sensor without cells.
Fig. 7. Proposed model for an electrode-solution-cell model with area A, uncovered by cells (a) and covered over area $A_c$ (b).

The relative changes in impedance magnitude, defined in eq. (12), inform more accurately from these variations: $r$ is the change in impedance magnitude for the two-electrode system with cells ($Z_c$) as compared to the system without cells ($Z_{nc}$). Figure 9 shows the $r$ versus frequency graph plotted for cell-to-electrode coverage $ff$ from 0.1 to 0.9 in steps of 0.1, using a $R_{gap} = 90 \, k\Omega$. The size of the electrode is $32 \times 32 \, \mu m^2$. Again we can identify the frequency range where the sensitivity to cells is high $[10 \, kHz, 1 \, MHz]$, represented by $r$ increments. For a given frequency, each value of the normalized impedance $r$ can be linked with its $ff$, allowing cell detection and estimation of the covered area $A_c$. Electrodes can be manufactured (in terms of technology and size) according to $r$-sensitivity curves to improve sensor response.

Fig. 8. Impedance evolution when fill factor increases $32 \times 32 \, \mu m^2$ for a $32 \times 32 \, \mu m^2$ electrode.

Fig. 9. Normalized impedance $r$ parameter evolution versus frequency. Each line corresponds to values of the fill factor (or $A_c$ covered by cells) in the range of 0.1 (nearly empty) to 0.9 (nearly full).
From Fig. 9, it can be deduced that electrode-cell electrical performance models can be used to derive the overlapping area in a cell-electrode system, an useful resource for biological studies. It can be observed how the curve matches with the frequency range, placing the maximum r value at around 100 kHz, as predicted in the COMSOL simulations. A value of $R_{\text{gap}} = 90 \, \text{k}\Omega$ was selected for this curve, representing a maximum value of the r curve with $\text{ff} = 0.69$, which in turn represents the ratio $(A_c/A)$, for the cell size of 30 $\mu$m diameter shown in the figure obtained using FEM simulations. Other cell sizes, with different diameters, can be selected from the r curves in Fig. 9.

6. AHDL model for simulation

Electrical simulations are commonly employed by electrical engineering designers to obtain useful information and predictions about circuit performance before circuits are sent to the factory. Simulations of front-end sensor data acquisition circuits must be trained versus real sensor models, so a reliable representation of a sensor is a prerequisite for a correct acquisition circuit design. The modelling process for an impedance sensor, as described in this work, can be performed using Analog Hardware Description Language (AHDL), which can easily be incorporated into mixed-mode simulators as SpectreHDL. Another advantage of using AHDL modelling is the possibility of incorporating non-linear circuit element performance, in our case the frequency square root function in the Warburg impedance. This makes AHDL models widely applicable, allowing accurate electrode-solution model simulation, and enabling their incorporation into mixed-mode simulators.

6.1 AHDL model

The proposed model has three main parameters: the electrode area ($A$), the fill-factor ($\text{ff}$) or percentage of the electrode area covered by cells, and the resistance of the gap region ($R_{\text{gap}}$). Technological data is also included, defining the physical properties of the sensor material and the solution. The HDL model directly implements the equations for circuit elements described in section 2: capacitor double layer, transfer resistance, Warburg resistance and spreading resistance, considering technology information (Borkholder, 1998) and parameters $A$, ff and $R_{\text{gap}}$. An example with the $R_{\text{ct}}$ resistance is described in the following lines:

```vhdl
module resistor_transfer(pin,nin) (A,FF,T) 
node [V,I] pin, nin;
parameter real A = 2500p;
parameter real FF = 0.0;
parameter real T = 309;
{
real curr_density = 2.0e-5;  /* Amp/m2 Au-sol reaction*/
real z_ions = 1.0;  /* Valence ions */
real q_e = 1.6e-19;  /* Charge e- */
real K = 1.3806e-23;  /* Constant Boltz. J/ºK */
real v_thermal = 0.0;  /* Thermal V. */
real rt = 0.0;  /* ohm.m2 */
analog {
v_thermal = K*T/q_e;
rt = (v_thermal/(curr_density*z_ions));
rt = rt / A*(1.0-FF);
```
Analog Hardware Description Language (AHDL) offers an easy way to describe equivalent circuits for the electrode-solution-cell model in Figure 1. Each circuit element is connected using its node description. Inside, the corresponding equations are incorporated to describe its electrical performance, technology data, geometry parameters (A) and model values (R_{\text{gap}}). This cell-electrode description can be used to carry out electrical simulations.

### 6.2 SpectreHDL simulations

In this section we present some simulation results obtained when modelling a commercial electrode: ECIS 8W10E, from Applied Biophysics Inc. (http://www.biophysics.com). A photograph of the complete system is shown in Fig. 10. Eight wells can be seen, each one containing ten circular gold microelectrodes, of 250 \( \mu \)m diameter. This impedance sensor will be employed for the cell culture experiments described in section 7. For electrical simulations, ten identical electrodes will be considered in parallel. The HDL model for each one is described using the circuit equations from section 2 and the model in Fig. 7b.

![Electrode employed for impedance measurements.](image)

Fig. 10. Electrode employed for impedance measurements.

The expected performance of ten circular electrodes is illustrated in the following simulations. Ten sensing electrodes were used (e_1 in Fig. 6), with just one common reference electrode (e_2), much larger than the sensing electrodes. Also shown are the electrode area dependence \( (A) \), the fill factor dependence \( (ff) \), and the resistance gap dependence \( (R_{\text{gap}}) \), to illustrate the model’s flexibility in different environments. Fig. 11 represents the normalized impedance r-values expected for these electrodes, for \( R_{\text{gap}} = 22 \text{ k}\Omega \), when the fill factor changes from a situation of electrodes uncovered by a cell \( (ff = 0.1) \) to nearly fully covered \( (ff = 0.9) \). Different cell densities or sizes can be detected using these sensitivity curves. Several values of \( R_{\text{gap}} \) can also be used to match the models to observed performance. In Fig. 12, several \( R_{\text{gap}} \) values were set for \( ff = 0.9 \), resulting in large changes in r-values. These parameters are used to apply the proposed model to a specific cell line or cell culture experiment. Finally, the electrode area for \( R_{\text{gap}} = 22 \text{ k}\Omega \) and \( ff = 0.9 \) was also changed. In this way, electrode frequency response can be set for the experiment’s given electrode geometry and technology parameters (Fig. 13). From these simulated results, it can be observed that optimal working frequency is close to the frequency proposed by the electrode factory (around 4 kHz), and that the electrode area covered by cells can be simulated using the fill factor parameters, making it possible to study cell density and size from the proposed model and simulations.
Fig. 11. The $r$ versus frequency curves obtained for $\eta \in [0.1,0.9]$ and $R_{\text{gap}} = 22 \, \text{k}\Omega$, using the data from 8W10E electrodes.

Fig. 12. The $r$ versus frequency curves obtained for $R_{\text{gap}} \in [10 \, \text{k}\Omega, 100 \, \text{k}\Omega]$ in steps of $10\,\text{k}\Omega$, for $\eta = 0.9$, using the data from electrodes 8W10E.

Fig. 13. The $r$ versus frequency curves obtained for $R_{\text{gap}} = 22 \, \text{k}\Omega$ and $\eta = 0.9$, for different values of the electrode area: 49n ($49.10^{-9} \, \text{m}^2$) corresponds to a circular electrode with a 250 $\mu\text{m}$ diameter.
6.3 Model extraction

The performance curves obtained above can be used to match experimental results to model parameters and obtain relevant biometric characteristics like cell size, cell culture density and gap resistance. Specific cell parameters such as size and cell density can be obtained from the fill factor values. Gap resistance value provides information about the cell type, because it is related to cell adhesion and cell membrane electrical performance. Models can thus be predicted from FEM simulations performed using the modelling procedure mentioned in section 2. This means that the experimental performance observed in labs can be matched to fit proposed model for the real time monitoring and interpretation of biological experiments, especially cell culture protocols.

7. Cell culture applications

In this section, we present some applications for the cell-electrode electrical model developed earlier. The objective is to employ the information obtained from model electrical simulations to interpret cell culture experiments and to provide information about cell size, cell culture growth processes and cell dosimetric characterization.

7.1 Size definition

Cell size can be measured directly from the fill-factor parameter using the sensor curve sensitivities in Figures 11 to 13. There, maximum sensitivities to cell size (maximum r values) are obtained at several frequency ranges (10 kHz to 1 MHz in Fig. 9). Cell size testing should be performed after detection of the optimum frequency range. Another consideration is that the electrode size should be of the same order as the cell size to be detected using the maximum sensitivity delivered by the impedance sensor. This figure demonstrates that cells of different sizes can be detected at a given working frequency by measuring their corresponding normalized impedance value.

7.2 Cell growth

In normal conditions, a cell generates one replica (divides in two equal cells) in a cell cycle, specific to each cell line, which may take several hours. Impedance sensors can be used to measure the number of cells in a cell culture experiment, because the cells are attached to the bottom electrode when they are not dividing. As the number of cells increases, the total bottom area covered by cells increases too, increasing the number of cells placed on top of an electrode and so in time, raising the impedance measured from the electrodes. In normal protocols, growth rate is measured by a tedious process that requires seeding a number of dishes equal to the number of time points to be represented, and counting the number of cells in each dish everyday. Fig. 14a shows a growth curve we measured over seven days using 8W10E sensors in a similar setup to that described in (Giaever, 1986). An initial number of approximately 5000 AA8 cells from chinese hamsters was seeded. From this curve, impedance dynamic range is around 1220 Ω, from 380 Ω to 1600 Ω. Considering the initial cell number of 5000 cells very low, we can take the initial impedance as due to the no-cell impedance value (Z_{nc}). At t = 6000 min, the medium was changed, and the confluent phase was achieved at t = 8500 min, approximately. The maximum experimental value given in eq. (12) is around r = 3.1, as illustrated in Fig. 14b. If we consider that in our model the electrodes are almost fully covered by cells (ff = 0.9), the best matching value of R_{gap} is 22 kΩ. System response corresponds to the r-values illustrated in figures 11 and 12. From these curves, the fill factors at different times can be obtained. Table I summarizes the
relative normalized impedance values $r$ at several times. Using Fig. 15 to represent the sensor response, it is possible to obtain the corresponding fill factor at every instant. The number of cells in real time can be estimated by considering that the number on top of each electrode can be measured in the order of 500 to 1000 cells. For a well area of 0.8 cm², the maximum number of cells goes from $0.8 \times 10^6$ to $1.6 \times 10^6$. The number of cells, $n_{	ext{cell}}$ in Table I, is obtained from $0.8 \times 10^6$, the expected final cell value. Initially, 5000 cells from the AA8 cell line were seeded. It was considered that a maximum 90% ($ff=0.9$) of the well surface was covered by the cells. A value of $Z_{nc} = 380\,\Omega$ for $r$ calculus at eq. (12) was also considered.

Fig. 14. (a) Impedance evolution in the cell growth experiment. (b) Estimated normalised impedance $r$ evolution.

![Fig. 14](image1)

Table I: Number of cells ($n_{	ext{cell}}$) obtained from the impedance $Z_c$ measure in Fig. 14a, and using sensor curves proposed for SW10E electrodes.

| time  | $r$   | $ff$ | $n_{	ext{cell}}$ |
|-------|-------|------|------------------|
| 0     | 0     | -    | 5000             |
| 500   | 0.024 | 0.020| 18000            |
| 1000  | 0.026 | 0.026| 44000            |
| 1200  | 0.072 | 0.070| 63000            |
| 2000  | n.a.  | n.a. | n.a.             |
| 2500  | n.a.  | n.a. | n.a.             |
| 3000  | 0.374 | 0.362| 322000           |
| 3500  | 0.437 | 0.395| 351000           |
| 4000  | 0.485 | 0.472| 447000           |
| 4500  | 0.510 | 0.506| 518000           |
| 5000  | 0.503 | 0.504| 535000           |
| 5500  | 0.500 | 0.500| 551000           |
| 6000  | 1.074 | 0.836| 551000           |
| 6500  | 1.587 | 0.784| 633000           |
| 7000  | 1.970 | 0.737| 686000           |
| 7500  | 2.333 | 0.810| 726000           |
| 8000  | 2.837 | 0.864| 768000           |
| 9000  | 3.234 | 0.906| 800000           |
| 9500  | 3.361 | 0.875| 780000           |
| 10000 | 3.500 | 0.843| 800000           |

Fig. 15. Table I shows the normalized impedance $r$-values obtained experimentally from sensor curves in Fig. 14. $ff$ values in Table I are obtained from experimental $r$-values, using curves in this figure. For example, at $t = 7500\,\text{min}$, the measured $r$ value is 2.353, which corresponds to $ff = 0.810$ on the impedance sensor curves, at the frequency of 4 kHz, and using $R_{\text{gap}} = 22\,\text{k}\Omega$. 

![Fig. 15](image2)
The approximate cell number can be deduced from the last column, allowing a real time cell count based on the r-curves associated with the sensors. If different types of cells are considered, for each one the $R_{gap}$ parameter can be calculated, producing r-curves specific to each cell line. This process allows cell lines to be identified and differentiated, for example to match the impedance measure of normal cells and cancer cells (Aberg, 2004).

### 7.3 Dosimetry analysis

Experiments to characterize the influence of certain drugs on cell growth were also developed. These are usually known as dosimetry protocols, and consist on determining the response of the cell growth to several drug doses. The objective is to demonstrate that the proposed model for a cell-electrode system allows cells to be selected and counted under different conditions. In our case, we again considered the AA8 cell line and, as drugs, we used six different doses of MG132 for growth inhibition, from 0.2 μM to 50 μM. After 72 hours of normal cell growth, the medium was changed and different doses of the drug were: 0.2, 0.5, 1, 5, 10 and 50 μM for wells 3 to 8 respectively. Well 2 was the control.

The experimental impedances obtained for the 8 wells are shown in Fig. 17, for a working frequency of 4 kHz. At the end of the experiment it could be observed that impedance decreases as drug dosage increases. The control well (W2) was full of cells with the maximum impedance, while the well with the maximum dosage (W8) had the lowest resistance, at the bottom. The black line (W1) represents electrode-solution impedance. In this case, after the medium change ($t = 4000$ min), impedance was seen inexplicably to drop below the initial baseline level (400 Ω). As in cell cultivation experiments, a resistance value was taken based on a starting value of $Z_{nc}$, as a value representative of electrode-solution impedance. Final impedance values at 8000 min, $Z_c$, were considered as the final response, in Table II. Considering $Z_{nc}$ and $Z_c$, r values are calculated in the third column. The r versus frequency curves in Fig. 15, can be used to obtain the estimated ff values from the proposed model. The number of cells at the end of the experiment was also counted, and is shown in the last column for each well. Considering $ff_{max} = 0.9$ for an experimentally measured cell number of $8.06 \times 10^5$, it is possible to calculate the rest of the expected ff values.

The same data is summarized in Table III for 2, 4 and 10 kHz respectively. The best match is obtained at 4 kHz in fill factor (ff parameter). The impedance baseline, $Z_{nc}$, for r calculus can be seen to decrease with frequency, due probably to electrode impedance dependence. For medium resistance (W1) and high drug concentration wells (W7, W8), the resistance measured is below $Z_{nc}$, so eq. (12) cannot be applied for r calculation.

### 8. Cell imaging impedance-based

Another application of the proposed model for testing cell culture impedance sensing is the implementation of an impedance based imaging system (Yúfera, 2011)(Yúfera, 2010a). This can be done when, as was commented before, cell sizes are of the same order as microelectrode sensors. In this example, we have only considered a simplified model for electrode-solution-cell electrical performance.

#### 8.1 Cell location applications

The cell-electrode model: Fig. 19 shows the model employed for a two-electrode sensor as shown in Fig. 6. For the empty electrode, the impedance model $Z(\omega)$ has been chosen. The
Fig. 16. Impedance obtained in 8 wells (W1 to W8) at 2 kHz frequency. W1: Medium. W2: Control. W3: 0.2 µM. W4: 0.5 µM. W5: 1 µM. W6: 5 µM. W7: 10 µM. W8: 50 µM.

Fig. 17. Impedance obtained in 8 wells at 4 kHz frequency. (a) W1: Medium (b) W2: Control (c) W3: 0.2 µM (d) W4: 0.5 µM (e) W5: 1 µM (f) W6: 5 µM (g) W7: 10 µM and (h) W8: 50 µM.

Fig. 18. Impedance obtained in 8 wells at 10 kHz frequency. (a) W1: Medium (b) W2: Control (c) W3: 0.2 µM (d) W4: 0.5 µM (e) W5: 1 µM (f) W6: 5 µM (g) W7: 10 µM and (h) W8: 50 µM.
Table II. Experimental values for relative impedance (r) and fill-factor (ff). $Z_{nc} = 400 \Omega$.

| Well | $Z_c$ at t= 8000 min | r from $Z_c$ and $Z_{nc}$ | ff estimated from model | ff expected | Nº° cells measured |
|------|----------------------|-----------------------------|-------------------------|-------------|-------------------|
| 1    | 259.2                | -                           | -                       | -           | -                 |
| 2    | 1631.7               | 3.1                         | 0.90                    | 0.900       | 8.06x10^8         |
| 3    | 1454.7               | 2.6                         | 0.85                    | 0.690       | 6.13x10^8         |
| 4    | 1030.6               | 1.5                         | 0.72                    | 0.610       | 5.41x10^8         |
| 5    | 625.8                | 0.5                         | 0.44                    | 0.410       | 3.60x10^8         |
| 6    | 417.4                | 0.05                        | 0.037                   | 0.036       | 3.20x10^4         |
| 7    | 406.8                | 0.015                       | 0.016                   | 0.024       | 2.10x10^4         |
| 8    | 99.6                 | < 0                         | -                       | 0.005       | 4.00x10^3         |

Table III. Experimental values for relative impedance (r) and fill-factor (ff) at different frequencies. $Z_{nc} = 480\Omega$, 400\Omega, and 315\Omega$ for 2, 4 and 10 kHz working frequency respectively.

| Well | $Z_c$ [\Omega] at t= 8000 min | r from $Z_c$ and $Z_{nc}$ | ff Estimated from model | ff expected | Nº° cells measured |
|------|--------------------------------|---------------------------|-------------------------|-------------|-------------------|
| 1    | 212                            | -                         | -                       | -           | -                 |
| 2    | 1500                           | 2.43                      | 3.1                     | 0.98        | 0.900             |
| 3    | 1524                           | 2.18                      | 2.6                     | 0.94        | 0.850             |
| 4    | 1042                           | 1.17                      | 1.5                     | 0.82        | 0.820             |
| 5    | 580                            | 0.21                      | 0.5                     | 0.32        | 0.440             |
| 6    | 432                            | -                         | 0.05                    | -           | 0.037             |
| 7    | 419                            | -                         | 0.015                   | -           | 0.016             |
| 8    | 84                             | -                         | 0.996                   | -           | 0.005             |

8.2 2D image applications

The case simulated was that of an 8x8 two-electrode array. The sample input to be analysed was a low density MCF-7 epithelial breast cancer cell culture, as shown in Fig. 20a. In this image some areas are covered by cells and others are empty. Our objective was to use the area parametrized electrode-cell model and the proposed circuits to detect their location. The selected pixel size was 50 \(\mu\)m x 50 \(\mu\)m, similar to cell dimensions. Figure 20a shows the grid selected and its overlap with the image. We associated a squared impedance sensor,
Fig. 19. Electrical models for (a) $e_1$ electrode without cells and, (b) $e_1$ cell-electrode. (c) Model for $Z(\omega)$. (d) Normalized magnitude impedance $r$ for $ff$ = 0.1 to 0.9 in steps of 0.1. $C_p = 1 \text{nF}$, $R_p = 1 \text{M}\Omega$, $R_s = 1 \text{k}\Omega$ and $R_{\text{gap}} = 100 \text{k}\Omega$.

Similar to the one described in Fig. 6, with each pixel in Fig. 20a to obtain a 2D system description valid for electrical simulations. An optimum pixel size can be obtained using design curves for normalized impedance $r$ and its frequency dependence. Each electrical circuit associated to each $e_1$ electrode in the array was initialized with its corresponding fill factor ($ff$) producing the matrix in Fig 20b. Each electrode or pixel was associated with a number in the range $[0,1]$ ($ff$) depending on its overlap with cells on top. These numbers were calculated to an accuracy of 0.05 from the image in Fig 20(a). The $ff$ matrix represents the input of the system to be simulated. Electrical simulations of the full system were performed at 10 kHz to obtain the impedance corresponding to each electrode using the AHDL model proposed in section 6. Pixels were simulated by rows, from the bottom left hand corner (pixel 1) to the top right hand corner (pixel 64) (Yúfera, 2011).

To have a 2D graph image of the fill factor (the area covered by cells) in all pixels, values or $r$ were obtained from the measurements taken and the sensor curves are used in Fig. 19. The results are shown in Fig 21, which represents the 8x8 $ff$-maps. In the maps, each pixel has a

Fig. 20. (a) 8x8 pixel area selection in epithelial breast cancer cell culture. (b) Fill factor map ($ff$) associated with each electrode (pixel).
Fig. 21. 2D diagram of the fill factor maps for 8x8 pixels: (a) ideal input. Image reconstructed from simulations at (b) 10 kHz and (c) 100 kHz.

grey level depending on its fill factor value (white is empty and black full). Specifically, Fig. 21(a) shows the ff-map for the input image in Fig. 20b. Considering the parametrized curves in Fig. 19 at 10 kHz frequency, the fill factor parameter was calculated for each electrode and the results are shown in Fig. 21b. The same simulations were performed at 100 kHz, producing the ff-map in Fig. 21c. As Fig. 19 predicts, the best match with the input is found at 100 kHz, since normalized impedance is more sensitive and the sensor has a higher dynamic range at 100 kHz than at 10 kHz. For reported simulations, 160 ms and 16 ms per frame are required working at 10 kHz and 100 kHz, respectively. This frame acquisition time is enough for real time monitoring of cell culture systems.

9. Conclusions and discussion

The main objective of the work described in this chapter is to develop alternative methods for measuring and identifying cells involved in a variety of experiments, including cell cultures. To this end, we have focussed on obtaining models of the sensor system employed for data acquisition, and on using them to extract relevant information such as cell size, density, growth rate, dosimetry, etc.

First of all, the impedance parameter was selected as an excellent indicator of many biological processes, and Electrical Cell-substrate Impedance Spectroscopy (ECIS), a technique currently considered very promising, was employed. Microelectrode impedance based sensors were analyzed, and finite element simulations carried out to model the electrical performance of both electrode-solution and electrode-solution-cell systems. This modelling process, starting with basic descriptions such as cell size, impedance, etc., enabled the use of different cells.
A practical circuit for electrode-solution and electrode-solution-cell simulation was employed, with the capacity to incorporate modelling information derived from FEM simulations. An Analog Hardware Description Model was proposed to incorporate this model to mixed-mode simulations at SpectreHDL, allowing non-linear performance characterization for the model proposed and the use of mixed-mode simulators. Commercial 8W10E electrodes supplied by Applied Biophysics (AB) were modelled using the proposed five element circuit, obtaining good matches. Optimal measurement frequency was identified near 4 kHz, which is the frequency usually recommended by AB. Relative impedance changes were also related to the cell density thanks to fill factor parameters.

A set of experiments were conducted to match the proposed models with the performance observed at cell cultures. The first proposal was to study cell growth evolution based on 8W10E electrodes. Curves obtained experimentally, using a basic set-up, allowed real time growth to be monitored. An estimation of the number of cells was obtained using sensor curves calculated from the proposed electrical model.

Dosimetry experiments reproduce conditions similar to those of cell growth, but in this case, a growth inhibitor was added at different doses. In the experimental data, it was observed that the higher the dose, the more the measure impedance decreased, in according with the expected performance. We still do not fully understand the experimental results obtained from the measurements. Impedance inexplicably fell below the expected baseline ($Z_{nc}$). However, impedance curves for the control and for small drug doses were perfectly aligned up to this impedance level. A proposed model was set up with $R_{gap} = 22 \, \text{k}\Omega$ to explain the experimental data. Deviations from the data expected were in the order of 10% and 20% in fill factor calculations (ff), and were more precise at 4 kHz.

A final application was illustrated in the field of microscopy for cell cultures through the correct decoding of impedance response. By calculating the fill factor, it is possible to define the area occupied over one microelectrode.

As mentioned earlier, deviation in the fill factor determination was seen to be large, and the results are not accurate; in the future it will be necessary to analyze the influence of certain error sources to increase the system performance. Firstly, data was collected using a simple set-up, in an attempt to reproduce the procedures followed in the experiments done by Giaever and Keese (1986). In these measurements, serial load resistance was large (1 MΩ), to limit the current through the cell to 1 μA. Signal-to-Noise Ratio (SNR) was low due to the large noise contribution from high resistance values. Other, improved setups should be used in the future. Secondly, the proposed model has the advantage that it needs only one parameter ($R_{gap}$), in comparison with another reported model using three parameters (Giaever, 1991). This makes easy to handle experimental data, but also introduces inaccuracies. The possibility of adding more parameters to the model should be considered in the future. Finally, finite element simulations must be developed for different cell lines to really test the cell identity in biometric applications.

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The methods for human identity authentication based on biometrics—such as the physiological and behavioural characteristics of a person—have been evolving continuously and seen significant improvement in performance and robustness over the last few years. However, most of the systems reported perform well in controlled operating scenarios, and their performance deteriorates significantly under real world operating conditions, and far from satisfactory in terms of robustness and accuracy, vulnerability to fraud and forgery, and use of acceptable and appropriate authentication protocols. To address some challenges, and the requirements of new and emerging applications, and for seamless diffusion of biometrics in society, there is a need for development of novel paradigms and protocols, and improved algorithms and authentication techniques. This book volume on Advanced Biometric Technologies is dedicated to the work being pursued by researchers around the world in this area, and includes some of the recent findings and their applications to address the challenges and emerging requirements for biometric based identity authentication systems. The book consists of 18 Chapters and is divided into four sections namely novel approaches, advanced algorithms, emerging applications and the multimodal fusion. The book was reviewed by editors Dr. Girija Chetty and Dr. Jucheng Yang. We deeply appreciate the efforts of our guest editors: Dr. Norman Poh, Dr. Loris Nanni, Dr. Jianjiang Feng, Dr. Dongsun Park and Dr. Sook Yoon, as well as a number of anonymous reviewers.

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