Chemical Compositions of Field and Globular Cluster RR Lyrae Stars. II. ω Centauri*

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Abstract

We present a detailed spectroscopic analysis of RR Lyrae (RRL) variables in the globular cluster NGC 5139 (ω Cen). We collected optical (4580–5330 Å), high-resolution (R ~ 34,000), high signal-to-noise ratio (~200) spectra for 113 RRLs with the multifiber spectrograph M2FS at the Magellan/Clay Telescope at Las Campanas Observatory. We also analyzed high-resolution (R ~ 26,000) spectra for 122 RRLs collected with FLAMES/GIRAFFE at the Very Large Telescope, available in the ESO archive. The current sample doubles the literature findings also suggest that two different RRL populations coexist in the cluster. The former is more metal-poor (~0.5 dex), with σ = 0.33 dex. Our findings also suggest that two different RRL populations coexist in the cluster. The former is more metal-poor (~0.5 dex), with σ = 0.33 dex. The latter is less numerous, more metal-rich, and yttrium enhanced ([Y/Fe] ~ 0.4). This peculiar bimodal enrichment only shows up in the s-process element, and it is not observed among lighter elements, whose [X/Fe] ratios are typical for Galactic globular clusters.

Key words: globular clusters: individual (NGC 5139) – stars: abundances – stars: variables: RR Lyrae – techniques: spectroscopic

Supporting material: machine-readable tables

1. Introduction

ω Cen (NGC 5139) is the most massive cluster in the Galaxy (4.05 × 106 M☉; D’Souza & Rix 2013), containing ~1.7 × 106 stars (Castellani et al. 2007). ω Cen is known to host stars that cover a broad range in metallicity, from [Fe/H] ~ −2.5 to [Fe/H] ~ 0.0 (Calamida et al. 2009; Johnson & Pilachowski 2010; Marino et al. 2011; Pancino et al. 2011; Villanova et al. 2014). This large metallicity spread, coupled with an age spread of ~2 Gyr (Villanova et al. 2014), suggests that ω Cen should be identified as the remnant core of a larger pristine dwarf galaxy, successively accreted by the Milky Way (Bekki & Freeman 2003; Da Costa & Coleman 2008; Marconi et al. 2014; Ibata et al. 2019). On the other hand, many studies suggest different origins, with ω Cen as the possible result of successive merging of inhomogeneous, coeval, protocluster clouds (Tsujimoto & Shigeyama 2003), or the result of a self-enrichment history within the cluster itself (Cunha et al. 2002). A general consensus about this peculiar cluster has not been reached.

Despite the uncertainties about its origin, ω Cen has several advantages related to its peculiar characteristics. Its huge number of stars permits estimation of its distance with multiple techniques, such as variable stars like Miras (Feast 1965),...
SX Phoenicis (McNamara 2000), Type II Cepheids (Matsunaga et al. 2006), and RR Lyrae (RRls; Braga et al. 2018; Boni et al. 2019), the tip of the red giant branch (RGB; Boni et al. 2008), or the white dwarf cooling sequence (Calamida et al. 2008). Among them, the large population of candidate RRls (~200 stars; Navarrete et al. 2015; Braga et al. 2018), makes $\omega$ Cen the ideal laboratory for a large investigation with multi-object spectroscopy. The multiple possibilities for a distance estimates, the large metallicity spread, and the large number of stars, provide a unique possibility to calibrate RRL period–luminosity–metallicity (PLZ) and period–Wesenheit–metallicity relations with a high level of accuracy, which then can be applied to other RRL samples in the Galaxy.

Photometric investigations concerning the RRls in $\omega$ Cen date back to more than one century ago (Bailey 1902) and they have been crucial objects for understanding the pulsation and evolutionary properties of old, low-mass helium-burning variables (Martin & Plummer 1915; Baade 1958; Sandage 1981a, 1981b; Boni et al. 2001, 2003). Optical time series CCD data were collected both by OGLE (Udalski et al. 1992) and by CASE (Kaluzny et al. 2005) experiments, and more recently by Weldrake et al. (2007). More recently, a complete optical (UBVRI; Braga et al. 2016) and near-infrared (JHK$_s$; Navarrete et al. 2015; Braga et al. 2018) census has been published. As usual, the high-resolution spectroscopic investigations lag when compared to the photometric ones. Some abundance analyses have been performed on the $\omega$ Cen RRls, based either on spectroscopic (Gratton et al. 1986, 18 stars), on spectrophotometric (Rey et al. 2000, 131 stars), or on photometric (Boni et al. 2019, 170 stars) techniques. However, the only large investigation based on high-resolution spectroscopy was performed by Sollima et al. (2006; 74 stars collected at $R \sim 22,500$). This work aims at improving the sample of available high-resolution spectroscopic abundances for the RRls in $\omega$ Cen, based on the techniques already applied in Magurno et al. (2018, hereafter Paper I) for the smaller monometallic globular cluster NGC 3201.

We describe the collected data set and the instrument settings in Section 2. Section 3 describes the analysis of radial velocities. The investigation methodology is presented in Section 4, and the abundance results are shown in Section 5 for iron, in Section 6 for the $\alpha$-elements, in Section 7 for the iron-peak elements, and in Section 8 for the yttrium. Finally, conclusions are presented in Section 9.
restricted to those stars within the field of view of M2FS. The Clement et al. online database is not independent of the other two. We also have our own positions from the FourStar (Persson et al. 2013) data set, already used in Braga et al. (2018). These have high accuracy because the pixel scale of FourStar is 0′′16/pixel and the typical image FWHM for the infrared photometry is ∼0′′5 or better. We first checked the targets positions because the M2FS fibers are placed in predrilled holes that must be accurate. We started with a sample of all the relevant stars, weeded out the ones for which Clement et al. have doubts, used the FourStar images to delete crowded stars, and adopted the FourStar positions where appropriate. This sample contained 160 RRLs, with roughly equal numbers of RRab and RRc. The sample was divided into 113 RRLs and 10 non-RRL stars. Only one RRL, V38, was observed twice. The main body of the paper only refers to the RRLs, whereas the non-RRL stars are briefly described in the Appendix.

In addition to our M2FS data, we also analyzed a sample of 560 multi-epoch spectra for 122 RRLs from the ESO archive, collected with the multiple-object, medium–high-resolution spectrograph FLAMES/GIRAFFE (Pasquini et al. 2002). We selected from the archive all the available RRL spectra collected with the HR13 grism, covering the wavelength range 6120–6405 Å with a spectral resolution R ≈ 26,400.

In total, 22 RRLs were only observed with M2FS, 31 RRLs were only observed with GIRAFFE, and 91 RRLs have spectra collected with both M2FS and GIRAFFE.

3. Radial Velocities

Estimating the radial velocity (RV) is a common way to establish whether a star is a globular cluster member. ω Cen has had many RV membership investigations thanks to its huge stellar population. Recently, An et al. (2017) estimated a cluster average velocity of 232.7 ± 0.6 km s⁻¹, with a dispersion σ = 14.4 km s⁻¹, by using 581 RGB stars. A decade earlier, Reijns et al. (2006) performed the largest investigation of ω Cen, estimating an average radial velocity of 231.3 ± 0.3 km s⁻¹ (σ = 11.7 km s⁻¹), with 1589 RGB stars. This very large cluster RV makes it unlikely that a field star in its sightline could be erroneously identified as a cluster member.

Nevertheless, we are dealing with variable stars, and single-epoch measurements are affected by intrinsic radial velocity variations along the pulsation cycle. Indeed, RRL pulsation cycles cause variations up to ∼70 km s⁻¹ in the observed RVs for RRab and up to ∼45 km s⁻¹ for RRC. Therefore, it was necessary to correct their observed radial velocities for the pulsational components, in order to determine their systemic (cluster) velocities, applying the velocity templates described in the following.

We first measured the instantaneous radial velocities using the task fxcor in IRAF (Tody 1986, 1993). The individual spectra were cross-correlated with a synthetic spectrum generated with the driver synth of the local thermodynamic equilibrium (LTE) line analysis code MOOG (Sneden 1973). This model spectrum was computed with the atmospheric parameters typical of stars in the RRL domain (Teff = 6500 K, log g = 2.5, ξturb = 3.0 km s⁻¹, [Fe/H] = −1.5; For et al. 2011; Sneden et al. 2017) and then smoothed to the M2FS or GIRAFFE resolution. The individual velocities are listed in the last column of Table 1, and we assume an average error for the entire sample of ∼1.3 km s⁻¹, as given by fxcor.

The use of multiple RV measurements allows us to improve the phasing of the individual data. The phase of the individual measurements was computed by using the period and the epoch of maximum light, relying on the work by Braga et al. (2016, 2018) for the most updated and homogeneous photometry, in the UBVRIJHKs bands, of the ω Cen RRLs. However, this approach is prone to possible systematics in cases of a large time interval between photometric and spectroscopic observations. Indeed, small errors in the determination of the period and/or in the epoch of maximum light could transform into large errors in the phase determination. Note that typical RRL periods range over about

21 Based on observations collected at the European Southern Observatory under ESO programme 074.B-0170(A), 074.B-0170(B), 082.D-0424(A), 081.D-0255(A).

22 IRAF is distributed by the National Optical Astronomy Observatories, which are operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.

23 http://www.as.utexas.edu/~chris/moog.html
6–18 hr (0.25–0.75 days). Moreover, for RRLs that are located in the cluster outskirts, we still lack an accurate epoch of maximum light (Navarrete et al. 2015; Braga et al. 2018). On the other hand, radial velocities are measured with high precision and have no dependence on photometry. Therefore, we can use RVs to compute more precise phases of the individual data points. To do that, we first defined two radial velocity templates, for RRab and RRc stars. Sesar (2012) identified a linear relation between the photometric V-band amplitude ($A_{V}$, mag) and the RV pulsation amplitude ($A_{v}$, km s$^{-1}$) for the RRab stars. Thus, we adopted his Equation (2),

$$A_{v,RRab} = 25.6(\pm 2.5)A_{V} + 35.0(\pm 2.3),$$

(1)

to scale his radial velocity curve template at the specific amplitude of each RRab star in our sample. The same approach was applied to the RRc stars, from the photometry and RVs presented by Sneden et al. (2017). We used the data in their Table 1 to define an average ratio between the velocity amplitude and the photometric V-band amplitude,

$$A_{v,RRc} = 54.7(\pm 2.7)A_{V},$$

(2)

and we scaled their radial velocity curve template accordingly. The next step was to fix the relative phases between the multiple RV measurements for a single star, according to their epochs and to the period. Finally, we used a minimization procedure with two free parameters (phase and average template velocity) and two fixed ones (measured RV and template amplitude) to phase our data. Figure 3 shows the alignment of the measured RV points with the RV template curves after the minimization procedure, for an RRab (top panel) and an RRc (bottom panel) star. The higher the number of points, the higher the precision of the result. We applied this method to all the stars for which at least three RVs were available (113 stars), and we used the usual method of maximum light epoch for the remaining ones (31 stars), using the most updated epochs for the RRLs in ω Cen estimated by Braga et al. (2016, 2018).

The top panel of Figure 4 shows the dependence on phase of the instantaneous radial velocity for both the M2FS (red filled circles) and GIRAFFE (black open circles) data sets. Pulsational velocity effects are easily seen in this panel. For the individual stars, we applied the template velocity curves to remove these effects and derive the systemic velocities ($V_{s}$). For the stars with three or more RV measurement, we computed $V_{s}$ as the integral average of the fitting template computed before. For the other stars, we anchored the template curve, scaled to the appropriate amplitude, to our single measured radial velocity and phase, based on the epoch of maximum light, and we computed the integral average velocity on the template curve. The bottom panel of Figure 4 shows that the estimated $V_{s}$ is almost independent of phase, within the natural star-to-star scatter. The average cluster velocity, from the joint samples of M2FS and GIRAFFE instantaneous velocities, was estimated as 232.6 ± 0.7 km s$^{-1}$, with a dispersion $\sigma = 17.1$ km s$^{-1}$. Once the template is applied, the average cluster velocity based on $V_{s}$ is slightly reduced to 231.8 ± 0.5 km s$^{-1}$, with a dispersion $\sigma = 13.9$ km s$^{-1}$. This value is in very good agreement with the cluster velocities found by Reijns et al. (2006) and An et al. (2017).

Figure 3. Results of the minimization procedure to phase multiple radial velocity measurements. The template curves by Sesar (2012) (RRab; top panel) and by Sneden et al. (2017) (RRc; bottom panel) are used as a reference to phase GIRAFFE (open black circles) and M2FS (filled red circles) observations.

4. Abundance Analysis

We used an equivalent-width (EW) analysis method to derive atmospheric parameters, metallicities, and relative abundances from the M2FS sample spectra.

4.1. Methodology

We selected the 140 atomic transitions listed in Table 2 from a collection of laboratory measurements and reverse solar analysis. This set of lines includes all of the transitions used in Paper I (see their Table 3 and references therein), augmented by some other lines that are detectable in the more metal-rich RRLs of ω Cen. We measured the EWs of these lines by means of a multi-Gaussian fitting performed with the pyEW code developed by M. Adamow. 25 Highly asymmetric lines were discarded, as well as too weak (EW < 15 mÅ) or too strong (EW > 180 mÅ) lines. The measurement error $\epsilon$ on the EW, for each absorption line, can be estimated using the relation by Venn et al. (2012),

$$\epsilon = (S/N)^{-1} \times 1.5 \times \text{FWHM} \times \delta x + 0.1 \times \text{EW},$$

(3)

where $\delta x$ is the pixel size of the instrument (180 mÅ). We obtained an average error for the entire sample $\epsilon \approx 8$ mÅ. As a final step, we used the LTE line analysis code MOOG, implemented in the Python wrapper pymOOG, to estimate atmospheric parameters ($T_{\text{eff}}$, log $g$, $\xi_{\text{turb}}$, [Fe/H]) and some relative abundances, using models interpolated from a grid of $\alpha$-enhanced (+0.4 in the log) atmospheres.

24 https://github.com/madamow/pyEW
25 https://github.com/madamow/pymoogi
26 We adopted the standard notation, [X/H] = $A(X) - A_{\odot}(X)$, where $A(X) = \log(N(X) / N_{\odot}) - 12$. Solar abundances refer to Asplund et al. (2009) within the text.
Table 2
Line List and Atomic Parameters

| λ (Å) | Species | EP (eV) | log(gf) (dex) |
|-------|---------|---------|--------------|
| 4702.991 | Mg I | 4.346 | -0.44 |
| 5172.684 | Mg I | 2.712 | -0.39 |
| 5183.604 | Mg I | 2.717 | -0.17 |
| 5266.556 | Ca I | 2.523 | -0.26 |
| 5031.021 | Sc II | 1.357 | -0.40 |
| 5239.813 | Sc II | 1.456 | -0.77 |
| 4981.731 | Ti I | 0.848 | +0.57 |
| 4999.503 | Ti I | 0.825 | +0.32 |
| 5064.653 | Ti I | 0.048 | -0.94 |
| 5173.743 | Ti I | 0.000 | -1.06 |

Errors on Iron Abundances Associated with Errors on the Parameter Estimates

| Species | Δlog g (±0.5 dex) | ΔΔTeff (±500 K) | ΔΔξturb (±0.5 km s⁻¹) |
|---------|------------------|----------------|------------------|
| Fe I/H | ±0.04 | ±0.01 | ±0.04 |
| Fe II/H | ±0.17 | ±0.04 | ±0.07 |

4.2. Metallicity-scale Calibration

This study and Paper I represent the first use of M2FS, with its limited spectral coverage, in a traditional abundance analysis of RRL stars. It is important to understand how the metallicity scale from our analysis compares with previous studies. To accomplish this, we used spectroscopic data from the high-resolution study of field RRLs recently reported by Chadid et al. (2017, hereafter C17). They collected thousands of spectra for a sample of 35 field RRab stars, with the du Pont telescope at Las Campanas Observatory, over several years. Their spectra cover a very large spectral interval, in the range 3400–9000 Å, much larger than the included M2FS spectral range, with a spectral resolution R ≈ 27,000. We performed our analysis on a selection of 27 stacked spectra (S/N ≈ 100) by C17, by using only the selected iron lines in the M2FS spectral range. In Table 4, we list the model parameters from C17 and from our M2FS analysis, along with the offsets between the two metallicity estimates. The agreement in the parameter sets is excellent: defining ΔX ≡ X_{C17} − X_{M2FS}, we found ΔTeff = −43 K (σ = 157 K), Δ log g = −0.07 (σ = 0.29), Δξturb = −0.33 km s⁻¹ (σ = 0.43 km s⁻¹), and Δ[Fe/H] = −0.02 (σ = 0.11). Most importantly, our M2FS-based Fe abundances are in very good agreement with the values obtained from the more comprehensive spectra of C17 (see Figure 5). A small offset can be noticed only for two out of the three C17 most metal-rich spectra; however, the differences are within 3σ from the mean. The difference of the third spectrum is still within 1σ. We can conclude that we are working on the same metallicity scale.

An additional calibration was performed in Paper I, in which the same kind of analysis, applied to the RRLs in the monometallic globular cluster NGC 3201, gave comparable results with previous studies based on nonvariable red giant stars.

4.3. Stellar Parameters

A total of 58 M2FS ω Cen spectra (57 objects) showed enough useful lines to perform a full spectroscopic parameter determination and abundance analysis. Figure 6 compares the relation Teff−log g for our M2FS sample (filled red circles), with the parameters obtained by For et al. (2011) and Sneden et al. (2017) for field RRLs (open black marks). The agreement of the two samples is good, with a few exceptions. In particular, two stars (V91 and V125) appear cooler than the bulk of the data. However, a visual inspection of the spectra does not give

References. Mg I, NIST database (Kramida et al. 2018), Ca I, NIST, Sc II, NIST, Ti I (Lawler et al. 2013), Ti II (Wood et al. 2013), Cr I (Sobeck et al. 2007), Cr II (Lawler et al. 2017), Fe I (O’Brien et al. 1991; Den Hartog et al. 2014; Ruffoni et al. 2014; Belmonte et al. 2017), Fe II, NIST, Ni I (Wood et al. 2014), Zn I, VALD database (Ryabchikova et al. 2015), Y II (Biémont et al. 2011).

(This table is available in its entirety in machine-readable form.)

Castelli & Kurucz 2003. The effective temperature was estimated by minimizing the dependence of the abundances on the excitation potential (EP), for the individual Fe I lines. The surface gravity was estimated by forcing the balance between the neutral and the ionized iron line abundances. Finally, the microturbulence was estimated by minimizing the dependence of the abundances on the reduced EW, RW ≡ log(EW/λ), for the individual Fe I lines.

Errors in estimating the atmospheric parameters also are reflected in the estimated abundances. Table 3 shows the
any argument to reject these stars as non-RRLs, so they are kept in the sample.

Another 51 M2FS spectra did not have enough iron lines to retrieve reliable atmospheric parameters with the EW method. The S/N of the spectra is quite homogeneous. The lack of lines is caused either by the low metallicity of the target, or to a hotter pulsation phase, or both. In particular, many of them did not have enough measurable Fe I lines to estimate effective temperature from Boltzmann excitation equilibrium, and others did not have any Fe II lines to estimate surface gravity from Saha ionization equilibrium. However, we were able to estimate average parameters starting from their phase. For et al. (2011) analyzed 11 field RRab, covering their entire pulsation cycles with multiple observations, showing that the atmospheric parameters have a relatively slow and regular variation along the pulsation cycle. The same applies to the 19 RRC analyzed by Sneden et al. (2017). However, the two quoted groups show, at fixed pulsation phase, a significant difference in the spread for which we do not have yet an explanation (see Figure 7). Both the samples were collected at the du Pont telescope and were analyzed with the same approach adopted by C17. This guarantees that we are still in the same calibration system as shown in Section 4.2. We applied the PERiodic GAuSsian Uniform and Smooth fit (PEGASUS) procedure described by Inno et al. (2015) to fit the atmospheric parameter distributions as a function of phase (solid lines in Figure 7). This was applied to the two individual samples of RRab and RRC, to obtain phase average parameters (hereafter called PAP) to be used in the abundance determinations. The fitting function is in the form

\[ y(\phi) = A_0 + \sum_{i=1}^{N} A_i e^{-B_i \sin^2(\phi - \phi_i)} , \]

where \( y \) is one of the atmospheric parameters (\( T_{\text{eff}}, \log g, \xi_{\text{surf}} \)) and \( \phi \) is the pulsation phase. All the coefficients are provided in Table 5. Unfortunately, the errors based on this approach are about one order of magnitude larger than those based on an EW analysis, for two reasons. (i) The atmospheric parameters are the average ones, and their standard deviations can be as high as \( \sigma_{T_{\text{eff}}} \approx 400 \text{ K}, \sigma_{\log g} \approx 0.6 \text{ dex}, \sigma_{\xi_{\text{surf}}} \approx 0.6 \text{ km s}^{-1} \), especially for the first overtone mode and during the phases of maximum light. This causes uncertainties in the abundances up to 0.4–0.5 dex (see Table 3). (ii) The spectra are not good candidates for a full EW analysis due to the paucity of good lines. This means that it is more difficult to decide whether a line is good or not with respect to the others, simply because there are few lines to compare with. Indeed, in a group of tens of lines, an outlier is immediately identified and removed. At the contrary, with only one or two lines it is not possible to exclude any value.

However, this approach gives better results than an estimate of the parameters based on photometric colors, as used, for example, by Sollima et al. (2006) and Johnson & Pilachowski (2010). To confirm that, we applied both the PAP and the
photometric approach to the sample of RRLs for which we spectroscopically estimated the atmospheric parameters. For the photometric approach, we used the parameterizations defined by Johnson & Pilachowski (2010), with the light curves in $V$ and $K_s$ collected by Braga et al. (2016, 2018). The microturbulence was defined by minimizing the abundance dependence on the reduced EW, once fixed $T_{\text{eff}}$ and $\log g$. The average differences, in terms of the atmospheric parameters, between the spectroscopic and the photometric estimates, confirm that the PAP appears to be more accurate. Indeed, defining $\Delta X_{\text{PAP}} = X_{\text{PAP}} - X_{\text{EW}}$ and $\Delta X_{\text{photometric}} = X_{\text{photometric}} - X_{\text{EW}}$, where $X$ represents one of the atmospheric parameters, we found $\langle \Delta T_{\text{eff,PE}} \rangle \approx 240$ K, $\langle \Delta T_{\text{eff,phot}} \rangle \approx 370$ K, $\langle \Delta \log g_{\text{PE}} \rangle \approx 0.04$, $\langle \Delta \log g_{\text{phot}} \rangle \approx 0.6$, $\langle \Delta \xi_{\text{turb,PE}} \rangle \approx 0.5$ km s$^{-1}$, $\langle \Delta \xi_{\text{turb,phot}} \rangle \approx 0.08$ km s$^{-1}$. The parameter dispersions in the two approaches are similar, of the order of $\sigma_{\text{PE}} \approx 500$ K, $\sigma_{\log g} \approx 0.6$, $\sigma_{\xi_{\text{turb}}} \approx 0.5$ km s$^{-1}$. We then applied the PAP approach to retrieve the abundances for the additional 51 RRLs.

For the remaining four spectra, no abundance analysis was possible because of the absence of useful lines or of phase information.

5. Metallicity Distribution

The iron abundance estimates for the individual stars are listed in Table 6. The sample of 57 RRLs for which we applied a full spectroscopic analysis based on the EW method shows an average cluster metallicity $[\text{Fe}/\text{H}] = -1.76 \pm 0.05$ and a large star-to-star dispersion $\sigma = 0.36$, as expected for $\omega$ Cen (Freeman & Rodgers 1975; Pancino et al. 2000; Calamida et al. 2009; Bono et al. 2019). As shown in Figure 8, top panel, our sample has its metallicity peak at about $[\text{Fe}/\text{H}] = -1.9$ and a pronounced tail toward higher metallicities, up to $[\text{Fe}/\text{H}] = -0.85$. The low-metallicity tail is much less evident, with the most metal-poor RRL estimated at $[\text{Fe}/\text{H}] = -2.53$.

Before taking into account the 51 additional RRLs obtained with the PAP approach, we performed a further calibration by computing, for the RRLs in the EW sample, the corresponding iron abundances with the PAP method. In Figure 9, we plotted the difference in $[\text{Fe}/\text{H}]$ between the two approaches, for the same stars, as a function of the iron abundance estimated with...
the PAP approach. There is clearly a large spread in the points, because the average atmospheric parameters can have higher or lower values than the real ones. Moreover, the parameterization for the RRab appears very promising, with a difference between the two approaches very close to zero, whereas the RRc appear, on average, more metal-rich with the PAP approximation. We therefore applied a zero-point calibration to the PAP sample of \( \omega \) Cen RRLs, according to the pulsation type, to make it consistent with the more accurate spectroscopic one. We also applied the same kind of correction to all the other elements, after performing a similar calibration based on their \([X/H]\) abundances. Figure 8, middle panel, shows the histogram for the entire M2FS sample (black thick line) after the calibration, together with the two subsamples: the EW (orange filled area) and the PAP (purple shaded area). For the joint EW and PAP samples, we derived \( \langle [\text{Fe/H}] \rangle \approx -1.82 \pm 0.03 \) (\( \delta = 0.33 \)). This mean value is only 0.06 dex lower than that derived with the pure EW analysis. Once again, the distribution peaks at about \( [\text{Fe/H}] \approx -1.9 \), with a longer metal-rich tail and a shorter metal-poor one.

Table 7 shows the mean iron content of the two RRL populations, RRab and RRc, with the different approaches adopted. It can be noticed that the EW method produces very similar results for both RRab and RRc, with a difference in the average iron content limited to 0.03 dex. On the contrary, the PAP method produces an RRab population that is 0.13 dex more metal-rich than the RRc one. In particular, the RRab sample has the same average abundance \( \langle [\text{Fe/H}] \rangle = -1.78 \)
Among the \(~500\) GIRAFFE spectra for which we measured a radial velocity, only a limited sample of \(~99\) spectra (44 objects, 27 in common with the M2FS sample) showed high enough \(S/N\) \((40 \leq S/N \leq 110)\) and useful iron lines to perform accurate EW measurements. However, the number of iron lines was too limited for a spectroscopic determination of the atmospheric parameters. In particular, they lacked useful Fe II lines to balance the surface gravity. Therefore, we applied the PAP approach to estimate the atmospheric parameters, then the abundances, for all the stars with available phase information and good enough iron lines.

As a first step, we averaged the abundances for the stars with multiple GIRAFFE measurements. Then, we compared the stars in common between the GIRAFFE and the M2FS samples. This defined a zero-point calibration for RRab and RRc, used to move the entire GIRAFFE sample to the M2FS, spectroscopic, metallicity scale. After the scaling, we performed a weighted averaged of the abundances for the stars with multiple measurements of the two spectrographs (last column of Table 6), assuming the inverse square of the error as weight. We ended with a sample of 125 RRLs, whose distribution is shown in the bottom panel of Figure 8. The shape of the distribution still remains essentially the same, with \((\langle [\text{Fe/H}] \rangle = -1.80 \pm 0.03\) and a dispersion \(\sigma = 0.33\) dex.

### 5.2. Comparison with the Literature

The large dispersion in the metallicity of \(\omega\) Cen is a well-known attribute that has been investigated for decades. Previous studies of both RGBs (Norris & Da Costa 1995; Johnson & Pilachowski 2010; An et al. 2017; Mucciarelli et al. 2018, 2019) and RRLs (Butler et al. 1978; Gratton et al. 1986; Butler et al. 1996, 1997) showed metallicity spreads between 0.20 and 0.45 dex. In Figure 10, we collected the \([\text{Fe/H}]\) distributions for the largest and most recent studies. With the exception of Rey et al. (2000, hereafter R00), whose distribution is essentially symmetric, the histograms show the larger metal-rich tail distribution also found in the present study. However, the RRL-based analysis of Sollima et al. (2006, hereafter S06) and Bono et al. (2019, hereafter B19), as well as the RGB analysis by Mucciarelli et al. (2018, 2019), show the metal-rich tail as a separated secondary peak, whereas the current sample shows either a metal-rich secondary peak for \([\text{Fe/H}] \geq -1.5\) (M2FS EW) or a well–defined metal-rich shoulder (M2FS+GIRAFFE). This was already noticed, among the others, by Norris et al. (1996, 1997) with Ca abundance and kinematics data, but we will discuss this point in more detail in Section 8.

It is worth mentioning that the iron distribution peak in the literature is, on average, \(~0.2\) dex more metal-rich than our estimate, with the exception of B19, who found a slightly more metal-poor peak. These differences are mainly due to the techniques used to estimate the atmospheric parameters. Indeed, S06, Johnson & Pilachowski (2010), and An et al. (2017) used photometrically estimated parameters that, as already mentioned in Section 4.3, give slightly higher \(T_{\text{eff}}\) and \(\log g\). Comparing our GIRAFFE analysis with the sample by S06, which is included in our sample, we found an average difference, for the stars in
common, $\Delta[\text{Fe}/\text{H}]_{\text{GIRAFFE-S06}} = -0.13$. Since this difference cannot be due to the spectra, it can only be related to the applied technique. Finally, R00 used the photometric $hk$ index to indirectly estimate the metallicity, while B19 used a technique based on PLZ theoretical predictions. The comparison of the literature results with our entire $\omega$ Cen sample (M2FS+GIRAFFE) is shown in Figure 11, only considering the stars in common among each work and this one. The average differences in $[\text{Fe}/\text{H}]$ are $\langle \Delta[\text{Fe}/\text{H}] \rangle_{\text{R00}} = -0.24 \; (n = 89, \sigma = 0.31)$,

\begin{table}[h]
\centering
\caption{Mean Iron Abundance and Standard Deviation for the Analyzed Samples, with Distinction between Different RRL Pulsation Modes}
\begin{tabular}{lccc}
\hline
Mode & M2FS$_{\text{REW}}$ & & M2FS$_{\text{PAP}}$ & & GIRAFFE & \\
& $N$ & $[\langle \text{Fe}/\text{H} \rangle]$ & $\sigma$ & $N$ & $[\langle \text{Fe}/\text{H} \rangle]$ & $\sigma$ & $N$ & $[\langle \text{Fe}/\text{H} \rangle]$ & $\sigma$ \\
\hline
RRab & 28 & $-1.78$ & 0.34 & 14 & $-1.78$ & 0.26 & 37 & $-1.68$ & 0.28 \\
RRc & 29 & $-1.75$ & 0.39 & 37 & $-1.91$ & 0.31 & 7 & $-1.83$ & 0.26 \\
\hline
\end{tabular}
\end{table}

Note. The number of stars for each subsample is also indicated.

Figure 10. Metallicity distribution for the sample of RRLs (black; left column) and RGBs (purple; right column) in $\omega$ Cen, available in the literature (R00: Rey et al. 2000; S06: Sollima et al. 2006; B19: Bono et al. 2019; N95: Norris & Da Costa 1995; J10: Johnson & Pilachowski 2010; A17: An et al. 2017; M18: Mucciarelli et al. 2018, 2019). J10, A17, and M18 have been scaled to a maximum height of 35 for plotting reasons. They should be multiplied by $\sim 5$, $\sim 4$, and $\sim 1.5$ respectively, to obtain the real scale. The number of stars in the sample, the mean, and the standard deviation are labeled in the top right corner.

Figure 11. Comparison of the metallicity distribution for our entire $\omega$ Cen sample (M2FS+GIRAFFE; black) and the RRL samples available in the literature (R00: Rey et al. 2000; S06: Sollima et al. 2006; B19: Bono et al. 2019). For each panel, only the stars in common ($n$) between the two works are considered. Mean and standard deviation of the differences among the two samples ($\Delta[\text{Fe}/\text{H}] = \text{this work} - \text{literature}$) are labeled in the top right corners of each panel.

$\langle \Delta[\text{Fe}/\text{H}] \rangle_{\text{this work-S06}} = -0.16 \; (n = 65, \sigma = 0.31)$, and $\langle \Delta[\text{Fe}/\text{H}] \rangle_{\text{this work-B19}} = +0.20 \; (n = 120, \sigma = 0.39)$.

6. The $\alpha$-elements: Mg, Ca, and Ti

The M2FS spectra cover a relatively short wavelength range, limiting the $\alpha$-element line measurements to only three species: Mg, Ca, and Ti. Titanium is not a “pure” $\alpha$-element (its dominant isotope is $^{48}$Ti instead of $^{44}$Ti); however, its abundance at low metallicity usually mimics those of the other $\alpha$-elements. Moreover, titanium lines are the most numerous after iron in the M2FS wavelength range, and in some cases they are the only observable ones among the $\alpha$. Indeed, up to 13 Ti I and Ti II lines

Figure 10.

Figure 11.
were measured in a single spectrum, whereas Mg I lines were limited to three at most, and only a single Ca I line was measured, if any. We estimated their average cluster abundances for the EW sample as \( [\mathrm{Mg/Fe}] = 0.41 \pm 0.03, [\mathrm{Ca/Fe}] = 0.46 \pm 0.03, \) and \( [\mathrm{Ti/Fe}] = 0.44 \pm 0.03. \) The dispersion of Mg is the largest one \( (\sigma = 0.22), \) whereas Ca has the smallest dispersion \( (\sigma = 0.13), \) but also the lowest number of measurements, and Ti lies in between \( (\sigma = 0.19). \) Adding the RRLs analyzed with the PAP approach does not change significantly the final results, but they double the number of stars: \( [\mathrm{Mg/Fe}] = 0.43 \pm 0.03, [\mathrm{Ca/Fe}] = 0.47 \pm 0.03, \) and \( [\mathrm{Ti/Fe}] = 0.44 \pm 0.02, \) with exactly the same dispersions as before.

Table 8 lists the individual star abundances, and Figure 12 compares them with those of the field halo RRLs, collected with high-resolution \( (R \geq 25,000) \) spectroscopy, available in the literature (Clementini et al. 1995; Fernley & Barnes 1996; Lambert et al. 1996; Kolenberg et al. 2010; For et al. 2011; Hansen et al. 2011; Liu et al. 2013; Govea et al. 2014; Pancino et al. 2015; Chadid et al. 2017; Sneden et al. 2017; more details about the literature samples can be found in the Appendix of Paper I). The agreement of the two considered samples is evident. The running mean of the two groups is the same, within the dispersion, in the metallicity range covered by the \( \omega \) Cen stars, where the \( \alpha\)-element abundances are almost constant or slightly decreasing toward higher metallicities. We also computed the \( [\alpha/Fe] \) abundance for the individual RRLs as the biweight mean of the three considered element abundances (bottom panel of Figure 12). Details about this robust iterative estimator of location can be found in Beers et al. (1990). The \( [\alpha/Fe] \) abundance was only estimated for those RRLs showing lines of all the three elements. This limits the sample to 18 RRLs, but the homogeneity of the results is preserved. The cluster average abundance was estimated as \( [\alpha/Fe] = 0.41 \pm 0.02 \) \( (\sigma = 0.10). \)

Figure 13 compares our \( \alpha\)-element abundances with those derived in earlier investigations of \( \omega \) Cen red giants, and there is a general agreement. It can be noticed that our estimates of Mg abundances are more scattered than those by Norris & da Costa (1995), but with similar mean values. However, this difference can simply be caused by the different sample size \( (78 \text{ versus } 40 \text{ measurements}). \) On the contrary, Ca and Ti have similar dispersions between RRLs and RGBs, but with higher

![Figure 12](image)

**Figure 12**: \( \alpha\)-elements vs. iron abundances for the RRLs in \( \omega \) Cen. The stars analyzed with the EW approach are marked with red filled circles, those analyzed with the PAP approach are marked with magenta open circles. The black symbols mark the field halo RRLs collected with high-resolution spectroscopy from the literature. In each panel, the black error bar in the bottom left corner shows the mean individual errors for the literature sample.

abundances for our sample, especially for Ti, with respect to both Norris & da Costa (1995) and Johnson & Pichachowski (2010). The two different populations display quite similar
Figure 13. $\alpha$-elements distribution for RGB stars in $\omega$ Cen (N95: Norris & Da Costa 1995; J10: Johnson & Pilachowski 2010), compared with our RRLs sample. J10 has been scaled to a maximum value of 28 for plotting reasons.

Figure 14. Top panel: $\alpha$-elements (Mg+Ca+Ti) vs. iron abundances for the individual RRLs in $\omega$ Cen (red filled circles). The sample is compared with Galactic globular cluster (filled black diamonds; Pritzl et al. 2005; Carretta et al. 2009, 2010) and field halo RRLs (open purple diamonds; same samples as in Figure 12). The solid blue line shows the log-normal fit of the two joint samples, with the 1σ dispersion shown by the blue bar in the top left corner. Bottom panel: as in the top panel, but compared with field halo giants (black dots; Frebel 2010) and RHB–BHB field stars (orange–blue squares; For & Sneden 2010). The dashed blue line shows an extrapolation of the fit toward lower iron abundances.

$\alpha$-element abundances over the entire metallicity range. To further investigate the chemical enrichment of the $\alpha$-elements in different stellar components, we also compared our results for $\omega$ Cen with similar abundances available in the literature for Galactic globular clusters (Pritzl et al. 2005; Carretta et al. 2009, 2010), field halo red/blue horizontal branch stars (RHB, BHB; For & Sneden 2010), and kinematically selected field halo red giants (Frebel 2010). Figure 14 shows all the previous samples and a log-normal analytical fit of their [α/Fe] versus [Fe/H], computed with Equation (1) in Paper I. It is remarkable that all the observed components, which are RRLs or nonvariables, and which are cluster or field stars, agree with the fit within 1σ. This suggests that all these components experienced very similar chemical enrichment histories for these $\alpha$-elements, also supporting a common old ($t \geq 10$ Gyr) age for all of them.

Figure 15. As in Figure 12(b), but for abundances derived from the GIRAFFE spectra (blue filled circles).

From the GIRAFFE spectra, the only element other than Fe that was possible to measure with sufficient precision is Ca. We estimated [Ca/Fe] for 40 out of 44 stars in the sample, shown in Figure 15. Despite a few outliers, the average Ca abundance for the GIRAFFE spectra is in good agreement with the literature values for field halo RRLs. We estimated the average abundance for the GIRAFFE sample, excluding the evident outliers with a 1σ clipping procedure, as $\langle [\text{Ca}/\text{Fe}] \rangle = 0.42 \pm 0.03$, with a dispersion $\sigma = 0.17$.

7. The Iron-peak Elements: Sc, Cr, Ni, and Zn

We estimated abundances for a few heavier elements in the iron-peak group ($Z = 21–30$): Sc, Cr, Ni, and Zn, only observable in the M2FS spectra. There are significant differences in the number of useful lines among the four species; indeed, some of these elements are undetectable in many of the stars (Table 8). Cr is the most broadly represented element, observed in about 50% of the RRLs (≈70% for the EW sample alone), with up to 10 lines in the best case and at least a couple of lines for the majority of the spectra (either Cr I or Cr II). On the other hand, Zn I was only observed in a handful of stars, especially in the metal-rich tail of the sample, with only one or two lines. Sc II and Ni I both have few observed lines in the M2FS spectral range, but the number of RRLs showing them is between those with Cr and Zn. The average cluster abundances were estimated as $\langle [\text{Sc}/\text{Fe}] \rangle = 0.11 \pm 0.04$, $\langle [\text{Cr}/\text{Fe}] \rangle = 0.09 \pm 0.02$, $\langle [\text{Ni}/\text{Fe}] \rangle = 0.06 \pm 0.04$, and $\langle [\text{Zn}/\text{Fe}] \rangle = 0.30 \pm 0.05$.

Figure 16 shows the comparison between the abundances of the iron-peak elements for the individual $\omega$ Cen RRLs and for the field halo RRLs available in the literature (see the Appendix in Paper I for more details about the literature sample). The agreement of the two groups is very good. The running means of the two samples with metallicity are nearly the same. The dispersions, in the metallicity range covered by $\omega$ Cen, are also very similar between the two groups and to those of the $\alpha$-elements, for Sc ($\sigma_{\omega \text{ Cen}} = 0.21$, $\sigma_{\text{halo}} = 0.15$), Cr ($\sigma_{\omega \text{ Cen}} = 0.18$, $\sigma_{\text{halo}} = 0.09$), Ni ($\sigma_{\omega \text{ Cen}} = 0.17$, $\sigma_{\text{halo}} = 0.26$), and Zn ($\sigma_{\omega \text{ Cen}} = 0.11$, $\sigma_{\text{halo}} = 0.14$). Once again, the chemical enrichment history of the RRLs in the Galactic halo appears to be similar for both field and cluster stars.

Our Fe-group abundances from RRL stars are also in good agreement with those derived from the RGB samples. The results obtained by Norris & Da Costa (1995) and Johnson & Pilachowski (2010) for the $\omega$ Cen RGBs (Figure 17) are in general agreement with our RRL sample, with only limited
differences in the average values, suggesting similar enrichment histories for the two stellar groups.

8. The s-process Element: Y

Among the neutron-capture elements (Z > 30), in the M2FS spectral range, only Y II lines are easily observable in RRLs (Table 8). A couple of La II transitions are also present, but they are too weak to produce reliably detectable lines. Indeed, synthetic spectrum tests show that these lines are not observable for [La/H] < −0.8, even with very high S/N. Y II lines were measured in the entire metallicity range covered by ω Cen, with an average abundance for the entire sample ([Y/Fe]) = 0.25 ± 0.05 and a dispersion σ = 0.31. Figure 18 shows the comparison between the Y abundances of RRLs in ω Cen and the field halo RRLs (top panel; see the Appendix in Paper I for details) and RGs (bottom panel; Frebel et al. 2010). Two groups of stars in ω Cen show different levels of Y enhancement: about half of the sample is in very good agreement with the field stars, having about solar Y abundances; the other half of the RRLs shows a clear over-enhancement of Y, with [Y/Fe] > 0.4. Figure 19 shows the comparison of two spectra: a Y-enhanced RRL (V112; [Y/Fe] = +0.61, black line) and an almost solar one (V20; [Y/Fe] = −0.10, red line). The two stars are both RRab, observed at similar phase (ϕ = 0.28 and 0.36, respectively), and with similar iron abundance ([Fe/H] = −1.78 and −1.76, respectively). The two shown iron lines are, indeed, almost identical, whereas the yttrium lines are largely different one from each other.

Figure 16. Iron-peak elements vs. iron abundances for the RRLs in ω Cen. The stars analyzed with the EW approach are marked with red filled circles, those analyzed with the PAP approach are marked with magenta open circles. The black symbols mark the field halo RRLs collected with high-resolution spectroscopy from the literature. In each panel, the black error bar in the bottom left corner shows the mean individual errors for the literature sample.

Figure 17. Iron-peak and s-process elements distribution for RGB stars in ω Cen (N95: Norris & Da Costa 1995; J10: Johnson & Pilachowski 2010), compared with our RRL sample. Zn is not shown because it is missing in the RGB samples. J10 has been scaled to a maximum value of 28 for plotting reasons.

Figure 18. Top panel: yttrium vs. iron abundances for the RRLs in ω Cen. The stars analyzed with the EW approach are marked with red filled circles, those analyzed with the PAP approach are marked with magenta open circles. The black symbols mark the field halo RRLs collected with high-resolution spectroscopy from the literature. The black error bar in the bottom left corner shows the mean individual errors for the literature sample. Bottom panel: as in the top panel, but compared with field halo RGs (black dots; Frebel 2010).
The Y-enhanced group of RRLs appears to be mostly the metal-rich one in ω Cen ([Fe/H] ≥ −1.5), suggesting differential enrichments for two groups of RRLs. A few other objects with similar strong Y enhancement are also observed at lower metallicity ([Fe/H] ∼ −1.8). However, they are a minor fraction of the RRLs more metal-poor than [Fe/H] ≤ −1.5. The abrupt increase in the s-process element abundances with increasing [Fe/H] was first observed by Lloyd Evans (1983) and later confirmed by François et al. (1988), Paltoglou & Norris (1989), and Vanture et al. (1994), not only for Y, but also for La, Zr, Ba, and Nd. Johnson & Pilachowski (2010) estimated the La abundance for ∼800 RGB stars, finding a clear separation between the most metal-poor stars, with almost zero enhancement ([Fe/H] ≤ −1.6, [La/Fe] = 0.0), and the most metal-rich, La enhanced ([Fe/H] ≥ −1.6, [La/Fe] ≥ 0.4). The hypothesis by Smith et al. (2000) and Cunha et al. (2002) is that two different populations coexist in ω Cen, whose enrichment history was strictly related to their capability to retain the products of the low-velocity ejecta of asymptotic giant branch stars (rich in s-process elements), allowing a heavy self-enrichment of the second, metal-rich, stellar generation, over timescales of the order of 1 Gyr. We note that the two groups of RRLs, the solar-enhanced and the overenhanced, show similar radial distributions from the cluster center and similar kinematic properties (radial velocity, radial velocity dispersion). However, more statistics are required before we can reach a firm conclusion.

A different scenario was advanced by Romano et al. (2007), who suggested that the self-enrichment scenario is not able to reproduce the metallicity distribution of ω Cen, and that the best hypothesis is that of ω Cen as the remnant of a dwarf spheroidal galaxy, evolved in isolation and then accreted by the Milky Way. In favor of the working hypothesis suggested by Romano et al., let us mention that the Y-enhanced RRLs appear to be an isolated group in terms of Fe and Y abundances, i.e., the current data do not suggest a steady increase in Y when moving from metal-poor to metal-rich RRLs. It is also worth mentioning that Norris et al. (1996) suggested, on the basis of a large sample of Ca abundances of ω Cen red giants, that ω Cen might be the merging of two different globular cluster. This kind of enrichment in s-process elements has never been observed in other Galactic RRLs, and indeed, field RRLs do not show similar Y overabundances.

Investigations of field stars and globular clusters in dwarf galaxies suggest that only a very limited difference with respect to the Milky Way exists for Y and other s-process element abundances, showing almost zero enhancement with respect to the Sun (Tolstoy et al. 2009 and references therein). Ba abundance in the Fornax dwarf galaxy represents a remarkable exception. Letarte et al. (2010) found [Ba/Fe] ≃ 0.7 for the investigated stars, more metal-rich than [Fe/H] ≃ −1.0. However, no similar enhancement was found for Y, and the problem is still open. Even if we cannot easily distinguish two separate populations from the [Fe/H] data, as was for S06 and B19, our [Y/Fe] abundances confirm that two distinct populations, one more metal-poor and the other more metal-rich than [Fe/H] ≃ −1.5, coexist in ω Cen. However, overenhanced RRLs are also observed at [Fe/H] ∼ −1.8, and the most metal-rich RRL in our sample shows almost solar Y abundance, so that the distinction between the two metallicity groups is not strict.

9. Conclusion and Final Remarks

We performed a large investigation of RR Lyrae stars in the globular cluster ω Cen, using high-resolution, high-S/N spectroscopy. We almost doubled the current sample of optical high-resolution spectroscopic abundances of RRLs, adding 109 cluster stars, observed with M2FS at the Magellan/Clay Telescope, to the ∼140 field halo stars available in the literature.

ω Cen was confirmed as a complex cluster, with a broad metallicity range and multiple populations. Indeed, the samples of proprietary M2FS data and archive GIRAFFE data allowed us to estimate ⟨[Fe/H]⟩ = −1.80 ± 0.03, with a high dispersion σ = 0.33. However, the average cluster metallicity alone is not sufficient to describe its complex nature. In agreement with previous investigations of various ω Cen samples, we found a nonsymmetric distribution of Fe, with a peak at [Fe/H] ≃ −1.85 and extended tails both in the metal-poor and especially in the metal-rich regime. The peak of the distribution is ∼0.2 dex more metal-poor than previous estimates for the cluster, with the exception of the work by Bono et al. (2019), who found an even more metal-poor distribution. The α- (Mg, Ca, and Ti) and iron-peak (Sc, Cr, Ni, and Zn) elements investigated show similar chemical enrichments to other known globular clusters and field stars of similar metallicity. In particular, the agreement was found not only with RRL stars, as the ones in our sample, but in general with variable and nonvariable field halo stars (RHB, BHB, and RGB stars), thus suggesting similar enrichment histories for all the analyzed old halo components. The α-elements are slightly enhanced, as expected for old stars, with [α/Fe] = 0.41 ± 0.02. The iron-peak elements show almost solar abundances, with the exception of Zn that appears slightly enhanced. On the contrary, the s-process element Y abundance shows peculiar characteristics, suggesting that two distinct populations coexist in the cluster, with the more metal-rich tail ([Fe/H] ≥ −1.5) dominated by stars with a strong enhancement of s-process elements, well represented by the average abundance of [Y/Fe] ≥ 0.4, and the more metal-poor stars with almost solar abundance. This overenhancement of the metal-rich population has no comparison in the field halo RRLs, appearing to be a peculiar characteristic of ω Cen.

The cluster radial velocity was estimated with the help of multi-epoch observations and template velocity curves to remove the phase-to-phase variability due to pulsation for the individual
observed. We finally estimated the average velocity of $\omega$ Cen as $231.8 \pm 0.5 \pm 13.9 \text{ km s}^{-1}$, in perfect agreement with literature results (Reijns et al. 2006; An et al. 2017).

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### Appendix

#### The Non-RRL Stars in $\omega$ Cen

For 10 stars in our M2FS sample, either the spectra are significantly different from those expected for an RRL or the EW analysis produced equilibrium atmospheric parameters ($T_{\text{eff}}, \log g, \xi_{\text{turb}}$) that are not typical of RRLs. However, their radial velocities confirm that these stars are actual members of $\omega$ Cen. The hypothesis is that the wrong stars were observed at the telescope due to the crowding of the $\omega$ Cen central region. Since we are not able to uniquely identify these stars within the cluster, we name them as UNK (unknown), followed by a sequential number corresponding to the RRL that was supposed to be observed (e.g., UNK15 was supposed to be the RRL V15 in $\omega$ Cen). We report in Table 9 a brief summary of their essential atmospheric parameters and abundances. As the nature for these objects is uncertain, we report our results on them only for completeness, but we would recommend further investigations/observations before using them for scientific purposes.

| ID    | $T_{\text{eff}}$ | $\log g$ | $\xi_{\text{turb}}$ | $\text{[Fe/H]}$ |
|-------|------------------|----------|---------------------|-----------------|
| UNK15 | 8200             | 4.30     | 4.50                | $-1.21 \pm 0.01$|
| UNK19 | 7100             | 4.00     | 5.10                | $-1.49 \pm 0.02$|
| UNK90 | 5700             | 3.60     | 2.00                | $-1.15 \pm 0.04$|
| UNK109| 7500             | 3.90     | 2.50                | $-0.92 \pm 0.04$|
| UNK114| 5600             | 4.10     | 2.00                | $-2.12 \pm 0.07$|
| UNK118| 6900             | 4.90     | 0.60                | $-0.59 \pm 0.03$|
| UNK143| 5800             | 3.00     | 3.00                | $-2.08 \pm 0.04$|
| UNK146| 5300             | 0.20     | 2.25                | $-2.21 \pm 0.01$|
| UNK267| 4900             | 0.90     | 2.60                | $-3.01 \pm 0.05$|
| UNK277| 5800             | 3.00     | 1.80                | $-1.40 \pm 0.05$|

| [Mg/Fe] | [Ca/Fe] | [Ti/Fe] | [Sc/Fe] | [Cr/Fe] | [Ni/Fe] | [Zn/Fe] | [Y/Fe] |
|---------|---------|---------|---------|---------|---------|---------|--------|
| 0.19    | ...     | 0.61 ± 0.04 | ...     | ...     | ...     | ...     | ...    |
| -0.41   | ...     | 1.09 ± 0.14 | ...     | 0.21 ± 0.10 | 0.03 ± 0.07 | 0.22    | 0.61 ± 0.10 |
| 0.16    | ...     | 0.66 ± 0.08 | 0.44 ± 0.23 | 0.74    | 0.45 ± 0.50 | 0.23 ± 0.26 | 0.73    | 0.78 ± 0.21 |
| ...     | 0.64    | 0.90 ± 0.17 | 0.30    | 0.09 ± 0.08 | ...     | ...     | ...    |
| -0.71   | ...     | 0.72 ± 0.28 | 0.39    | 0.58     | ...     | ...     | 0.48    |
| 0.21    | ...     | 0.89 ± 0.08 | ...     | -0.20    | ...     | -0.07 ± 0.10 |
| 1.72    | ...     | 0.22 ± 0.18 | ...     | 0.71 ± 0.04 | ...     | ...     | -0.12    |
| ...     | 1.48    | 0.59 ± 0.02 | ...     | 0.19 ± 0.08 | 0.09 ± 0.16 | -0.09    | 0.23    |
