The Effect of Alloying Elements on the Structural Stability, and Mechanical and Electronic Properties of Al₃Sc: A First-Principles Study

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Abstract: The first-principles methods, based on the density function theory, are performed to calculate the properties of pure and doped Al₃Sc. The structural stability, and mechanical and electronic properties of L₁₂-Al₃Sc₁₋ₓMₓ (M = Zr, Ti, Y, and Li) have been investigated. A negative formation enthalpy for L₁₂-Al₃Sc₁₋ₓMₓ indicated that all doped structures were stable, and Al₂₄Sc₆Zr₂ was found to be the most stable. The elastic constants, elastic moduli and Debye temperatures of Al₃Sc, with different doping elements and different doping concentrations, were calculated to explore the influences of doping on the mechanical properties and Debye temperatures of Al₃Sc. Furthermore, the calculated results suggested that both Al₂₄Sc₆Zr₂ and Al₂₄Sc₆Ti₂ could optimize the mechanical properties. Finally, the electronic properties based on the analyses of densities of states and electron density distributions, have been performed, to explain the underlying mechanisms for the structural and mechanical properties of the L₁₂-Al₃Sc₁₋ₓMₓ structures.

Keywords: ab-initio calculations; structural stability; elastic properties; electronic properties

1. Introduction

In the past decades, Al-Sc alloys have attracted great attention, due to higher strength and stiffness, compared with pure Al [1,2]. The addition of Sc can greatly improve the mechanical properties of Al alloy, since a large number of nanoparticles (Al₃Sc) are formed in the aging process [3–5], and Al₃Sc with a cubic L₁₂ structure has a small lattice mismatch with α-Al. However, the industrial application of Al₃Sc has been limited because of the high cost of Sc, and the low solubility of Sc in Al. As a result, it is necessary to find an element to replace Sc, which can improve the mechanical properties of the alloy and reduce its cost.

In recent years, researcher has shown that doping has a certain impact on the structural characteristics and mechanical properties of alloys, through experimental investigations [6–9] and theoretical calculations [10–12]. The behaviors and properties of the alloying elements in Al₃Sc, have been studied by experiments. For example, Fuller et al. studied the replacement of Sc by Zr in Al-Sc alloys, and found that the coarsening resistance was increased at higher temperatures [13]. Dalen et al. studied the effects of Ti additions on the structural and creep properties of Al-Sc alloys [14]. Seidman et al. indicated that the addition of Li has resulted in an increased driving force for precipitate nucleation [15]. Harada et al. compared the thermal expansion and creep properties of Al₂₃Sc and Al₃(Sc, Y) [16,17]. Up to now, the microstructure, compression, fracture behavior, elastic and optical properties of Al₃Sc have been widely researched [18,19] by the density functional theory (DFT) method. Furthermore, the structural, electronic, mechanical, and thermodynamic properties of Al₃Sc, under
different pressures and temperatures, have been calculated [20–22]. Moreover, the effect of transition metals on the structural stability of Al$_3$Sc$_{1-x}$M$_x$ was studied by the special quasi random structures method [23]. In addition, the properties of Al$_3$Sc$_{1-x}$M$_x$ (M=Zr, Ti) with different concentrates, were performed [24,25]. Nevertheless, the influences of the doping elements on the mechanical and electronic properties of the Al-Sc-M system have required further investigations.

In this work, we calculated the structural stability and mechanical properties of the L1$_2$-Al$_3$Sc$_{1-x}$M$_x$ (M = Zr, Ti, Y, and Li) structures with concentrations of 3.125 at.% and 6.25 at.%. Second, the mechanical properties and electronic properties of the L1$_2$-Al$_3$Sc$_{1-x}$TM$_x$ structures have been intensively discussed in relation to their electronic properties. This investigation could provide theoretical guidance to the application of Al-Sc-based alloys.

2. Materials and Methods

All calculations were performed on the basis of the density functional theory (DFT) with the Vienna Ab Initio Simulation Package (VASP) [26]. The pseudopotential in the reciprocal space was described by the projector-augmented wave (PAW) method [27]. The generalized gradient approximation (GGA), with the Perdew-Burke-Ernzerhof (PBE) [28] function was applied to describe the exchange-correlation potential. Both the k-space integral and plane-wave basis were chosen to ensure that the total energy was converged. The convergence criterion for the self-consistent field energy was set to be $5 \times 10^{-6}$ eV/atom. For the plane wave expansion, a kinetic cutoff energy of 500 eV was considered to be sufficient. The geometry optimization was terminated when the Hellman–Feynman force on each atom was smaller than the 0.003 eV/Åmm. The integral in the Brillouin zone was sampled by the Monkhorst-Pack method [29], with the k-point mesh of 21 $\times$ 21 $\times$ 21 for Al$_3$Sc and 11 $\times$ 11 $\times$ 11 for the 2 $\times$ 2 $\times$ 2 supercells. All calculations were carried out with the potentials for Al(3$s^2$3$p^1$), Li(1$s^2$2$s^1$), Sc(3$s^2$3$p^6$3$d^1$4$s^2$), Ti(3$s^2$3$p^6$3$d^2$4$s^2$), Zr(4$s^2$4$p^6$4$d^2$5$s^2$), and Y(4$s^2$4$p^6$4$d^1$5$s^2$), as the valence electrons. Overall, all calculations were operated in 0 K, with the equivalent hydrostatic pressure.

3. Results

3.1. Structural Stability

The Al$_3$Sc phase had a cubic structure (Space Group: Pm3m (No. 221)), which contained 3 Al atoms and 1 Sc atom, in a unit cell. Based on the Al$_3$Sc phase, a 2 $\times$ 2 $\times$ 2 supercell (Figure 1a) was constructed to investigate the effects of the doped element (M = Zr, Ti, Y, and Li). In addition, the effects of the doping concentrations (at.%) were also considered, which included 3.125% (Figure 1b) and 6.25% (Figure 1c). The structures of the supercells used in this work are shown in Figure 1, and the optimized structural parameters are listed in Table 1.

![Figure 1](https://example.com/figure1.png)

**Figure 1.** Structures of Al$_3$Sc (a) 2 $\times$ 2 $\times$ 2 supercell; doped with the alloying element (M = Sc, Zr, Ti, Y, Li) at different doping concentrations, (b) 3.125%, and (c) 6.25%. Blue, green, and pink balls represent (M = Zr/Ti/Y/Li), Sc, and Al atoms, accordingly.
was consistent with results from previous studies [23,30]. In addition, it was found that the stability of the 
structure could stably form at 0 K; all structures were stable. However, the ∆Hf values of the doped 
structures were greater than that of pure AlSc, except for the Al-Sc-Zr system. This meant that Zr 
doping could improve stability, while the Ti/Y/Li dopings might reduce the structural stability. This 
was consistent with results from previous studies [23,30]. In addition, it was found that the stability 
of Al2Sc6Zr2 was higher in comparison to Al2Sc7Zr. Except for the Zr addition, the higher concentration 
dopings should have caused a lower stability of the structure.

3.2. Elastic Properties

The elastic properties could provide necessary information on the resistance of the material to 
extrinsically applied stress. For the cubic crystals, there were three independent elastic constants (i.e., 
C11, C12, and C44). The calculated results of the elastic constants are listed in Table 2.

| Structures | C11 (GPa) | C12 (GPa) | C44 (GPa) | C11 − C12 (GPa) | C12 − C44 (GPa) |
|------------|-----------|-----------|-----------|-----------------|-----------------|
| Al2Sc      | 183.99    | 37.81     | 71.94     | 146.18          | −34.13          |
| Exp. [31]  | 183       | 46        | 68        | 137             | −22             |
| DFT [15]   | 180.67    | 40.62     | 72        | 140.05          | −31.38          |
| Al2Sc6Zr   | 181.2     | 40.89     | 69.44     | 140.31          | −28.55          |
| Al2Sc6Zr2  | 183.93    | 42.9      | 69.07     | 141.03          | −26.17          |
| Al2Sc7Ti   | 181.66    | 40.88     | 69.58     | 140.79          | −28.7           |
| Al2Sc6Ti2  | 186.41    | 43.325    | 69.26     | 143.09          | −25.94          |
| Al2Sc7Y    | 175.77    | 38.08     | 68.59     | 137.69          | −30.51          |
| Al2Sc6Y2   | 171.96    | 38.55     | 67.93     | 133.41          | −29.38          |
| Al2Sc7Li   | 167.18    | 38.05     | 69.69     | 129.37          | −31.64          |
| Al2Sc6Li2  | 163.42    | 39.34     | 68.41     | 124.08          | −29.07          |

Note: Exp. represents the reported "experimental results".
Notably, the results of Al$_3$Sc were close to the published DFT calculations [15] and experimental measurements [31], implying the reliability of the calculation results. The criterions for the mechanical stability of the cubic crystal [32] were estimated using Equation (2):

$$C_{11} - C_{12} > 0, \ C_{11} > 0, \ C_{44} > 0, \ C_{11} + 2C_{12} > 0$$ (2)

The calculated results showed that the elastic constants of all structures could satisfy the above stability conditions, indicating that all structures had mechanical stability.

The calculated elastic constants as a function of doping concentration for Al$_3$Sc are exhibited in Figure 2. The $C_{11}$ values of Al$_3$Sc, before and after doping, were larger than the other elastic constants (Figure 2a), suggesting that the axes compression resistances were stronger. The $C_{11}$ values in the Al$_{24}$Sc$_7$M system were obviously decreased, which proved that the axes pressure resistances were reduced. However, Al$_{24}$Sc$_6$M$_2$ have found to be quite complex. For example, Al$_{24}$Sc$_6$Y$_2$ and Al$_{24}$Sc$_6$Li$_2$ have smaller $C_{11}$ values, Al$_{24}$Sc$_6$Zr$_2$ had a similar $C_{11}$ value, and Al$_{24}$Sc$_6$Ti$_2$ possessed a larger $C_{11}$ value. On the other hand, both Al$_{24}$Sc$_7$M and Al$_{24}$Sc$_6$M$_2$ exhibited higher $C_{12}$ values (Figure 2b), suggesting that the Poisson effect was enhanced [25]. Additionally, the Al-Sc-M systems had shown a slightly reduced tendency for $C_{44}$ (Figure 2c), especially for the Al$_{24}$Sc$_6$M$_2$, indicating that the resistance to shear deformation was gently diminished.

![Figure 2](image.png)

**Figure 2.** The elastic constants $C_{ij}$ (GPa) of L1$_2$-Al$_3$Sc doped with the element M (M = Zr, Ti, Y, and Li) on the dependence of doping concentration: (a) $C_{11}$, (b) $C_{12}$, and (c) $C_{44}$.

The polycrystalline elastic moduli are the important performance parameters of engineering materials, such as bulk modulus ($B$), shear modulus ($G$), Young’s modulus ($E$), and Poisson’s ratio ($\nu$). The bulk modulus and shear modulus were determined using the Voigt-Reuss-Hill method [33]. For the cubic structure, the bulk modulus ($B$) and shear modulus ($G$) were calculated from Equations (3) and (4), respectively:

$$B = B_V = B_R = (C_{11} + 2C_{12}) / 3$$ (3)

$$G_V = (C_{11} - C_{12} + 3C_{44}) / 5$$

$$G_R = 5(C_{11} - C_{12})C_{44} / [3(C_{11} - C_{12}) + 4C_{44}]$$ (4)

$$G = (G_V + G_R) / 2$$
Young’s modulus (E) and Poisson’s ratio (ν) were determined by B and G, and their expressions are shown in Equations (5) and (6) [34]:

\[ E = \frac{9GB}{3B + G} \]  \hspace{1cm} (5)

\[ \nu = \frac{3B - 2G}{2(3B + G)} \]  \hspace{1cm} (6)

The results of the elastic moduli are listed in Table 3. It is well-known that B is a measure of the degree to which a material deforms under hydrostatic pressure [35]. G indicates the material’s resistance to shear strain [35]. E is a representation of the stiffness of the material [36].

**Table 3.** Elastic moduli B, G, E, ν, B/G, H, A and Θ_D for L1_2-Al_3Sc_M compounds, at the ground state.

| Structures | B (GPa) | G (GPa) | E (GPa) | ν | B/G | A | H (GPa) | Θ_D (K) |
|------------|---------|---------|---------|---|-----|---|---------|---------|
| Al_3Sc     | 86.53   | 72.4    | 169.84  | 0.173 | 1.195 | 0.98 | 15.78   | 607     |
| Exp. [31]  | 91.7    | 71.7    | 170.63  | 0.201 | 1.28 | 0.99 |         |         |
| DFT [13]   | 87.3    | 71.2    | 167.94  | 0.179 | 1.23 | 1.03 |         |         |
| Al_3Sc_Zr  | 87.66   | 69.73   | 165.34  | 0.186 | 1.257 | 0.99 | 14.59   | 596     |
| Al_3Sc_Zr  | 89.91   | 69.64   | 166.06  | 0.192 | 1.291 | 0.98 | 14.3    | 590     |
| Al_3Sc_Ti  | 87.81   | 69.91   | 165.74  | 0.185 | 1.256 | 0.98 | 14.68   | 610     |
| Al_3Sc_Ti  | 91.02   | 70.16   | 167.48  | 0.193 | 1.297 | 0.97 | 14.37   | 608     |
| Al_3Sc_Y  | 93.98   | 68.64   | 161.93  | 0.179 | 1.223 | 0.99 | 14.7    | 594     |
| Al_3Sc_Y  | 83.02   | 67.44   | 159.21  | 0.18  | 1.23 | 1.02 | 14.39   | 580     |
| Al_3Sc_Li  | 81.09   | 67.64   | 158.77  | 0.174 | 1.199 | 1.07 | 14.61   | 611     |
| Al_3Sc_Li  | 80.97   | 65.78   | 155.43  | 0.18  | 1.231 | 1.10 | 14.05   | 619     |

Notes: A is defined for elastic anisotropy, and Θ_D is for Debye temperature.

The influences of doped elements at different concentrations on the elastic modulus of Al_3Sc are displayed in Figure 3a–c. The values of B doped with Zr and Ti increased with a growing doping concentration, while the values of B doped with Y and Li decreased with an increasing doping concentration. When the doping concentration was constant, Ti doped Al-Sc-M compounds had the maximum B values, implying that the addition of Ti into Al_3Sc resulted in minimal deformation, under certain external pressure.

![Figure 3](image-url)
Furthermore, the doped Al Sc-M compounds displayed lower $G$ and $E$, meaning that the shear strain resistance and stiffness were reduced. In addition, the $E$ values of Al$_3$Sc doped with Zr and Ti were increased, when the doping concentrations increased from 3.125% to 6.25%, indicating that the higher concentration doping showed a better performance for the Al Sc-M (M = Zr/Ti).

The investigation of the stiffness could be completed by providing the microhardness parameter ($H$), given by the following relation [37]:

\[
H = \frac{(1 - 2\nu)}{6(1 + \nu)} E
\]  

Figure 3d exhibits the $H$ values of Al$_3$Sc with different doping elements and concentrations. It is clearly observed that the $H$ values of the doped Al$_3$Sc decreased with an increasing concentration. For example, Al$_{24}$Sc$_6$Ti$_2$ presented the maximum $H$ value. Although the $H$ values of the doped Al$_3$Sc were decreased, the reduced magnitude was small, compared with that of Al$_3$Sc. Considering the cost of Sc, the doped Al$_3$Sc was more valuable for industrial application.

The B/G ratio was used to distinguish the ductility and brittleness of compounds [38]. The greater value of B/G corresponded to a better ductility in the material. The results shown in Figure 4a indicate that the L1$_2$-Al$_3$Sc structure displayed a higher ductility, after doping. This conclusion could be proved by the Cauchy pressure ($C_{12}$–$C_{44}$) [39]. Furthermore, the B/G values for the doped Al$_3$Sc increased with the elevated doping concentration, which demonstrated that a higher doping concentration enhances the ductility of Al$_3$Sc. In addition, Poisson’s ratio ($\nu$) reflected the transverse deformation for the material [40]. It could be clearly observed that the changing trends of $\nu$ with doping elements and concentrations were similar to that of B/G (Figure 4). It is common knowledge that a material exhibits better ductility, when the Poisson’s ratio is large. The calculated results demonstrated that a better ductility of Al$_3$Sc doped with Zr/Ti was exhibited, compared to Y/Li.

![Figure 4. (a) B/G and (C$_{12}$ – C$_{44}$) and (b) $\nu$ of L1$_2$-Al$_3$Sc doped with element M (M = Zr, Ti, Y, and Li) as a function of the doping concentration.](image)

The elastic anisotropy has an important implication in engineering science, since it is highly correlated with the possibility to induce micro cracks in materials [41]. The expression of elastic anisotropy ($A$) is shown below:

\[
A = \frac{2C_{44}}{(C_{11} - C_{12})}
\]  

It is noted that the material is isotropic when $A$ is equal to 1. The degree to which the value of $A$ deviates from 1 represents the strength of the anisotropy of the material. The $A$ values of Al-Sc-M (M = Zr/Ti/Y/Li) were calculated, and the results are listed in Table 3. The results showed that Al$_3$Sc had exhibited anisotropic behavior, before and after doping. Moreover, the $A$ values for the doped Al$_3$Sc were more deviated from 1, when the doping concentrations varied from 3.125% to 6.25%. It is worth
noting that the deviation degree of Al-Sc-Li system was the largest. For example, Al_{24}Sc_{6}Li_{2} possessed the strongest anisotropy.

### 3.3. Debye Temperature

Debye temperature ($\Theta_{D}$) is a fundamental parameter for the material’s thermodynamic properties. It is correlated with many physical properties (i.e., specific heat, elastic constant, and melting temperature). The $\Theta_{D}$ value of a solid can usually be calculated from the sound velocity. $\Theta_{D}$ is related to the modulus of elasticity in Anderson’s model [42]. Then, $\Theta_{D}$ is defined as [43]:

$$
\Theta_{D} = \frac{h}{k} \left[ \frac{3n}{4\pi} \left( \frac{N_{A}\rho}{M} \right) \right]^{1/3} v_{m}
$$

(9)

$$
v_{m} = \left[ \frac{1}{3} \left( \frac{2}{v_{l}^3} + \frac{1}{v_{l}^3} \right) \right]^{-1/3}
$$

(10)

$$
v_{l} = \frac{G}{\rho}
$$

(11)

$$
v_{l} = \frac{3B + 4G}{3\rho}
$$

(12)

where $h$, $k$, $N_{A}$, $n$, $\rho$, $M$, $v_{m}$, $v_{l}$, $v_{l}$, and $v_{l}$ denote Planck’s, Boltzmann’s and Avogadro’s constants, total number of atoms, density, molecular weight, average sound velocity, transverse sound velocity, and longitudinal sound velocity, respectively.

In the Debye theory, $\Theta_{D}$ is the temperature of a crystal’s highest normal mode of vibration. That is, the highest temperature can be achieved due to a single normal vibration. It is well-known that a higher $\Theta_{D}$ corresponds to a better thermal conductivity of a material. Figure 5 describes $\Theta_{D}$ of Al$_{3}$Sc, at different doping elements and concentrations. The $\Theta_{D}$ values of the Al-Sc-Li structure increased with an increasing concentration. On the contrary, the $\Theta_{D}$ values of Al-Sc-M (M = Zr/Ti/Y/Li) decrease with the increasing concentration. For the Al-Sc-Ti structure, the $\Theta_{D}$ values were a little higher than Al$_{3}$Sc. However, the $\Theta_{D}$ had decreased slightly, when the concentration increased from 3.125% to 6.25%. The higher $\Theta_{D}$ values of the Al-Sc-Li structure indicated that their thermal conductivities were better, compared to other structures.

**Figure 5.** The $\Theta_{D}$ values of L1$_{2}$-Al$_{3}$Sc doped with element M (M = Zr, Ti, Y, and Li) as a function of the doping concentration.

Based on the above calculation results, it was noticeable that the species and concentrations of the doped elements had intensive impacts on the mechanical properties of Al$_{3}$Sc. In order to obtain high
performance compounds, it was important to select appropriate doped element and concentration. The mechanical properties of Al$_3$Sc before and after doping were intensively compared. It was demonstrated that doping Ti/Zr could better optimize the performances of the Al$_3$Sc structure. This conclusion was consistent with previous reports [14,17,23]. Furthermore, the mechanical properties of Al-Sc-Ti were slightly better than that of Al-Sc-Zr. This discrepancy with the result, measured through experiments, might have arisen from a different temperature. It is worth noting that the experimental measurement was generally operated at 573 K, whereas, this work was completed at 0 K. To be more important, the structure with a higher doping concentration (6.25%) had a higher performance over the lower doping concentration (3.125%).

3.4. Electronic Properties

To gain a better understanding of the doping effects at the electronic level, the total densities of states (TDOS) and partial densities of states (PDOS) were calculated in this work. The calculated TDOS are shown in Figure 6. The Fermi energy level ($E_f$) represented by a dotted line was set to zero. It could be clearly seen that the TDOS of the structure was not zero at the Fermi level, indicating that the structures had good metallic properties [2]. Meanwhile, the TDOS of Al$_{24}$Sc$_6$M$_2$ (M = Zr/Ti/Y/Li) were expensed in the energy scales.

![Figure 6. Total density of states (TDOSs) of the Al$_3$Sc and Al$_{24}$Sc$_6$M$_2$ (M = Sc, Zr, Ti, Li and Y) alloys.](image)

The structures of Al$_3$Sc, before and after doping, had a very wide pseudogap around $E_f$, implying that Al-Sc-M intermetallic compounds had strong covalent bonds, indicating that it had a good stability [44]. Moreover, the widths of the pseudogap for Al$_3$Sc and Al-Sc-M (M = Zr, Ti, Y, and Li) structures were 2.65 eV, 2.74 eV, 2.53 eV, 2.59 eV, and 2.23 eV, respectively. It was indeed seen that Al$_{24}$Sc$_6$Zr$_2$ was slightly wider than the pseudogap of Al$_3$Sc, while that of Al$_{24}$Sc$_6$M$_2$ (M = Ti/Y/Li) was slightly narrower. This proved that Al$_{24}$Sc$_6$Zr$_2$ had a stronger covalent bond, which was the most stable [13,14,23]. However, the stability of M (M = Ti/Y/Li) doped Al$_3$Sc was weakened. This was consistent with the calculation of the formation energy in Table 1.

The reason for this consequence, the Zr-$d$ orbital provided more valence electrons to hybridize with the Al-$p$ orbital than the Sc-$d$ orbital, while the Ti-$d$ orbital provided fewer valence electrons. At the same time, there was a stronger $d$-$d$ bond interaction between the Sc and Zr atoms, which effectively enhanced the ductility of the material [45]. For the Al$_{24}$Sc$_6$Li$_2$ structure, the Li atom replaced the Sc atom, and reduced the $p$-$d$ hybridization. The number of bonding electrons per atom in Al$_3$Sc and Al-Sc-M (M = Zr, Ti, Y, and Li) structures in the low energy region are shown in Figure 7, which were 2.94, 3.022, 2.994, 2.943, and 2.811, accordingly (energy range between $-12$ eV and Fermi levels). It is well-known that a higher number of bonding electrons implies an increased structural stability [34,46]. Thus, a stronger electron interaction should occur in Al$_{24}$Sc$_6$Zr$_2$, and Al$_{24}$Sc$_6$Zr$_2$ should have a larger
structural stability. It was considered that as Zr had more valence electrons, it resulted in stronger electron interactions between the Zr-\(d\) orbital and the Al-\(p\) orbital, as well as between (Sc, Zr)-\(d\) [45,47].

![Figure 7. Total density of states (TDOSs) and Partial density of states (PDOSs) of (a) Al\(_3\)Sc and (b–e) Al\(_{24}\)Sc\(_6\)M\(_2\) (M = Zr, Ti, Y, and Li) alloys. TDOSs for different structures and different elements are represented by purple lines. PDOSs for different orbitals are denoted by \(s\) in red, \(p\) in green and \(d\) in blue.](image)

To further illustrate the contribution of each atomic orbital to TDOS, the PDOS of each atom were calculated, as shown in Figure 7. The main bonding peaks of Al\(_{24}\)Sc\(_6\)M\(_2\) were predominantly derived from the Al-\(s\) and Al-\(p\) orbitals, in the energy range between −12 eV and −4 eV (Figure 7b–e), making the TDOS of Al\(_{24}\)Sc\(_6\)M\(_2\) almost coincident with the Al\(_3\)Sc.

It could be clearly observed that, from −4 eV to 5 eV, the TDOS was mainly contributed by the strong hybridization of the Al(Li)-\(p\) and Sc(Zr/Ti/Y)-\(d\) orbitals, and a small contribution of Al-\(s\) was also observed. Additionally, there was a large overlap in the entire energy range, leading to a strong \(pd\) hybridization. The pseudogap was generated by the hybridization of Al-\(p\) and M-\(d\) (M = Sc, Zr, Ti, Y).
In other words, there was a strong covalent bond in the Al-Sc-M (M = Zr, Ti, Y, and Li) structure. In addition, the PDOS of Zr/Ti/Y-d orbitals were different from the Sc-d orbital, suggesting that the doping elements should have taken effect on the TDOS. Below the Fermi level, the peaks of Zr-d orbital moved toward the lower energy level (Figure 7b), contributing to the enhancement of the bonding states of Al$_{24}$Sc$_6$Zr$_2$. Regardless, the Sc-d orbital of Al$_{24}$Sc$_6$Li$_2$ moved toward the higher level (Figure 7c). Moreover, the magnitudes of Li-s/p orbitals were tiny, compared to the other atomic orbitals. This implied that there was a weaker covalent bond in Al$_{24}$Sc$_6$Li$_2$ for the subdued hybridization between Li and Sc atoms.

For a deeper insight into the atomic bonding of the doped structures, the valence electron density distribution were also investigated. For example, the charge densities on the (100) and (110) planes, for each cell, are shown in Figure 8, in which the contour lines are plotted from 0.015 to 0.04 e/Å$^3$ with 0.0025 e/Å$^3$ interval.

**Figure 8.** The electron density contour plots (Unit: e/Å$^3$) on the (100) and (110) planes for L1$_2$-Al$_3$Sc (a,b), L1$_2$-Al$_{24}$Sc$_6$Zr$_2$ (c,d), L1$_2$-Al$_{24}$Sc$_6$Ti$_2$ (e,f), L1$_2$-Al$_{24}$Sc$_6$Y$_2$ (g,h), and L1$_2$-Al$_{24}$Sc$_6$Li$_2$ (i,j).

Figure 8a,b display the charge distribution of pure Al$_3$Sc on the (100) and (110) planes, respectively. It was clearly observed that the charge densities of the neighboring Al-Al, Al-Sc, and Sc-Sc had overlaps,
especially between the Al-Sc, indicating that there were strong covalent bonds in Al₃Sc. Moreover, these covalent bonds were mainly generated by the hybridization between Al-p and Sc-d orbitals. Compared with (110) plane, the charge distribution between the neighboring Al-Sc on the (100) surface was weak, which proved that the covalent bond on the (110) plane was stronger, leading to a brittle fracture. This was caused by the difference in the local symmetries between the (100) and (110) plane. This feature was consistent with previous reports [36]. The charge densities of Al₂₄Sc₆Zr₂ (100) and the (110) planes are shown in Figure 8c,d, respectively. The overlap of the charge density between the neighboring Al-Zr was increased, which meant that the covalent bond was enhanced. Moreover, Zr was slightly less electronegative than Sc. On this account, Al-Zr exhibited weaker ionic bond properties. This property could be analyzed by the Bader charge [48,49]. The calculation results showed that 0.38 electrons were transferred from Sc to Al in Al₃Sc. Nevertheless, the charge transferred from (Sc, Zr) to Al was reduced to 0.34, and the charge distribution on the Sc atom remained unchanged. However, the bonding difference between (100) and (110) planes was decreased, which was profitable for the improvement of the ductility of Al₂₄Sc₆Zr₂ [24]. The charge distributions of the Al₃Sc doped with Ti and Y are shown in Figure 8e–h, accordingly. In contrast to the Al, charge densities between Al-Ti, Y, Sc-Ti, and Sc-Sc were reduced, implying that the covalent bonds were weakened. The charge density distributions of Al₂₃Sc doped with Li are shown in Figure 8i,j. It was observed that the charge density overlap between Sc-Li decreased, indicating that Sc-Li exhibited a weaker covalent bond, which was mainly contributed by the hybridization of the Sc-d and Li-p states.

4. Conclusions

In order to explore the effects of the doped elements (M) on the mechanical properties of Al₃Sc, both, the structural stability and mechanical properties of Al₃Sc with different doping elements and concentrations, in combination with the influence of the higher doping concentration on the electronic properties of Al₃Sc were systematically investigated using the first-principles methods. Based on the results of this study, the following conclusions could be deduced. First, it was observed that the Al₃Sc structure could be stable after doping. For instance, Al-Sc-Zr had the highest stability, and Al₂₄Sc₆Zr₂ performed better on stability, compared to Al₂₄Sc₆Zr. On the other hand, Al-Sc-M (M = Ti, Y, Li) reduced the stability of Al₃Sc and Al₂₄Sc₆M₂ performed worse on stability, compared with Al₂₄Sc₆M. The calculated elastic constants of Al₃Sc, before and after doping, showed its mechanical stability. Moreover, the calculated B/G results revealed that the doped Al₃Sc with a higher concentration exhibited a better ductility, especially, when doped with Zr and Ti. It was noted that, the calculated results of the elastic modulus B, G, E, and ν suggested that both Al₂₄Sc₆Zr₂ and Al₂₄Sc₆Ti₂ displayed better mechanical properties. Additionally, the TDOS and PDOS analyses indicated that the doped Al₃Sc had a pseudogap and a strong covalent bonding, which was due to the strong pd state hybridization. Among them, the maximum pseudogap existed in Al₂₄Sc₆Zr₂, indicating its best stability. This conclusion was consistent with the calculated formation enthalpy. Ultimately, the obtained results could provide an important theoretical basis for a wide application of the Al-Sc alloy.

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