The formation of stars in clusters

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Abstract. Observations of the dust and gas around embedded stellar clusters reveal some of the processes involved in their formation and evolution. Large scale mass infall with rates $\dot{M} \simeq 4 \times 10^{-4} M_\odot \, \text{yr}^{-1}$ is found to be disrupted on small scales by protostellar outflows. Observations of the size and velocity dispersion of clusters suggest that protostellar migration from their birthplace begins at very early times and is a potentially useful evolutionary indicator.

1. Introduction

Massive stars form in the densest regions of the most massive molecular clouds and they do so in groups. To study massive star formation, therefore, requires that we learn about the processes by which stars form in clusters.

Much of what we have learnt about star formation, however, has come from studying nearby, isolated, low mass stars (e.g. Evans 1999). Protostellar evolution passes through four stages, Class 0 to III, that are defined on the basis of their observed spectral energy distribution and is physically related to the mass of the molecular envelope around the star (André, Ward-Thompson, & Barsony 2000). As a protostar grows through accretion while simultaneously dispersing the surrounding gas through outflows, the column density of the envelope around it decreases and the emergent spectral energy distribution shifts to shorter wavelengths from far-infrared (Class 0) to visible (Class III). The earliest stages are the shortest lived, with statistically based estimates of the lifetime of the Class 0 stage $\sim 5 \times 10^4$ yr, and Class I $\sim 10^5$ yr compared to Class II and III lifetimes $> 10^6$ yr (Mundy, Looney, & Welch 2000).

Just as the availability of infrared arrays facilitated the study of clusters of Class II sources in the 1990s (e.g., Lada 1992), so the arrival of millimeter and sub-millimeter wavelength bolometer arrays has made possible large scale studies of clusters at the earlier Class 0–I stages of protostellar evolution (e.g. Johnstone, this volume) We are now in a position to address questions about cluster formation and evolution with a level of detail that approaches the case of isolated star formation.

There are important differences between the formation of stars in clusters and in isolation. For instance, cluster forming clouds may initiate collapse through the dissipation of turbulent support (Nakano 1998) rather than the quasi-static gravitational contraction of a magnetically supported core (Shu, Adams, & Lizano 1987) and may form stars more efficiently (Lada 1992). Massive stars in particular, drastically alter their environment (e.g., Walborn, this
volume) and may induce further star formation through cloud compression (Elmegreen & Lada 1977). On the other hand, they may also be the principal agents of cloud destruction (Williams & McKee 1997).

To learn about the processes by which clouds gather the material to form clusters and the effect of newly formed stars back on the cloud requires observations of the intra-cluster gas and dust. Here I discuss radio observations of two highly embedded, Class 0 and I protostellar groups and present a potential new method for following cluster evolution at these early stages.

2. Millimeter wavelength observations of NGC2264

NGC2264 lies in the northern region of Monoceros, 760 pc distant, and contains about 300 near-infrared and 30 IRAS sources. The former are mostly Class II protostars and the latter Class I. Using the Heinrich Hertz Telescope on Mt. Graham, AZ we observed dust continuum emission at 870 µm and J = 3 − 2 lines of HCO\(^+\) and H\(^{13}\)CO\(^+\) toward three protostellar groups in NGC2264 to search for redder, younger sources and to follow the motions of the gas around them. Details of the observing procedure and data analysis can be found in Williams & Garland (2001).

The location of the observations is shown in Figure 1. IRS1 is a bright, compact group centered on an early B star situated just north of the Cone nebula. IRS2 and IRS3 are less luminous IRAS dust condensations lying ∼ 6′ and 11′, respectively, to the north.

Figure 1: Digital sky survey and IRAS 100 µm images of the NGC2264 cluster forming region. Color selected IRAS and MSX point sources (stars and circles respectively) and the boundaries of the HHT 870 µm bolometer maps are overlaid on each image.
2.1. Structure

Thermal dust emission at 870 µm toward IRS1 and IRS2 is shown in Figure 2. IRS1 is the bright, compact source to the south. Its elongated shape shows signs of substructure; Ward-Thompson et al. (2000) resolve five distinct condensations with higher resolution IRAM and JCMT observations. IRS2 lies to the north and is more fragmented with several low intensity peaks, suggestive of a slightly more evolved cluster in which the individual protostars are beginning to shed their circumstellar envelopes and migrate away from their birthplace. IRS3 was not detected in a short integration and is only included here in speculative comments toward the end of the paper.

![Figure 2: 870 µm continuum emission toward IRS1 and IRS2. Black contours are at 0.5, 1.0, 1.5, 2.0, 2.5 Jy beam\(^{-1}\), and white contours are at 3.5, 4.5, 5.5, 6.5 Jy beam\(^{-1}\). The boundary of the map is marked by the heavy solid line and the 21" beam size is indicated in the lower left corner.](image)

The integrated fluxes, for flux densities greater than 0.5 Jy beam\(^{-1}\), of the two clusters are similar, 94 Jy for IRS1 and 84 Jy for IRS2 indicating that they have similar dust masses and make a good comparative study. This implies total (gas + dust) masses of 790 M\(_\odot\) (IRS1) and 710 M\(_\odot\) (IRS2) for the envelopes around each protostellar group, for a mass opacity \(\kappa = 0.009\) cm\(^2\) g\(^{-1}\) and dust temperature \(T_d = 17\) K. The projected area of each group, 0.45 pc\(^2\) for IRS1 and 0.62 pc\(^2\) for IRS2, implies average column and volume densities, \(\langle N_{\text{H}_2}\rangle =\)
The high column densities toward the protostars imply $A_V \gtrsim 50$ and suggest they are indeed young (Class 0–I) objects. The average density of the envelopes traced by the continuum observations is similar to the critical density of HCO$^+$. Further, integrated line maps show a similar morphology to the dust (Williams & Garland 2001) and we therefore use the line data to investigate the dynamics of the gas around the clusters and the protostars within. Each group is discussed separately: the $29'' \simeq 0.1$ pc resolution of the line data is too coarse to investigate the dynamics around the individual protostars in IRS1 and we focus on global properties for this group; in the larger IRS2 group, individual objects are resolved and the gas motions around each one are analyzed.

2.2. Dynamics

Both the optically thin H$^{13}$CO$^+$ and optically thick HCO$^+$ $J = 3 - 2$ lines were observed. The former allows us to analyze the systemic velocity, linewidth, and abundance in the clumps around each cluster and the latter provides information on relative motions, i.e. infall and outflow, in the gas. Because of the limited space, I can only discuss the infall motions here and refer the reader to Williams & Garland (2001) for the H$^{13}$CO$^+$ velocity and abundance analysis.

The HCO$^+$ spectra generally show two peaks and a central dip at the mean velocity of H$^{13}$CO$^+$. Such self-absorbed profiles can be used to diagnose relative motions in the gas (e.g., Walker et al. 1986). Average spectra are shown toward IRS1 in Figure 3. The averages were computed over different regions of the clump corresponding to different 870 $\mu$m flux densities. This is similar to radial averaging but takes into account the non-circular shape of the clump. At low flux densities, corresponding to the outer regions, the average spectra are relatively weak and the optical depths are low. Nevertheless, the HCO$^+$ spectra are clearly asymmetric and show the red-shifted self-absorption characteristic of collapse. At higher flux densities, toward the clump center, the line intensities and optical depths increase and the central self-absorption dip becomes more prominent. However, the spectra also become more symmetric with increasing flux density suggesting a decrease in the net inflow.

Determining the infall speed profile requires detailed modeling. We ran the 1-d radiative transfer code, RATRAN, of Hogerheijde & van der Tak (2000) and used the continuum and H$^{13}$CO$^+$ observations to fix the size, temperature, density profile, abundance, mean velocity and dispersion. Only the relative velocity profile and $^{12}$C/$^{13}$C isotope ratio remain as variables in the modeled output HCO$^+$ line profiles.

Model spectra are shown as dotted lines in Figure 3 and reproduces the qualitative features in the HCO$^+$ spectra reasonably well given the tight constraints from the dust and H$^{13}$CO$^+$ observations. The model consists of a spherical shell at a radius $r_{in} = 0.32$ pc, with mass $400 \ M_\odot$, infalling at speed $v_{in} = 0.3$ km s$^{-1}$ onto an expanding inner layer moving outwards at $v_{exp} = 1.0$ km s$^{-1}$. The abundance of HCO$^+$ was constant and equal to $1.5 \times 10^{-9}$ relative to H$_2$, similar to values in other OB star forming regions (Bergin 1997). The inferred mass infall rate onto the cluster is,

$$\dot{M}_{in} = M_{in}v_{in}/r_{in} = 4 \times 10^{-4} \ M_\odot yr^{-1}.$$
This is one to two orders of magnitude greater than infall rates onto solar and intermediate mass protostars Zhou (1995) as expected for a global infall onto a group of such objects and is similar to the value expected for a gravitational flow, \( M = \sigma^3/G \), where \( \sigma = 1.0 \) km s\(^{-1}\) is the average velocity dispersion of the core and \( G \) is the gravitational constant (Shu 1977).

![Image](NGC2264 IRS1)

**Figure 3:** Global properties in IRS1. The left panel shows the continuum emission smoothed to the 29\(^\prime\) resolution of the line data. Contours are at 1,2,3,4,5 Jy beam\(^{-1}\), and spectra averaged within each contour interval are shown on the right hand side (HCO\(^+\) in light, H\(^{13}\)CO\(^+\) in dark and multiplied by 3 for clarity). At low flux densities/large radii, the HCO\(^+\) optical depth is low and the self-absorption effect is small but a clear asymmetry can be seen indicating collapse. At high flux densities/small radii, the self-absorption dip is strong but the spectra are more symmetric indicating a smaller net infall rate. Model fits, as described in the text, are shown as dotted lines.

The expanding inner region is likely due to the protostellar outflows in the region, one of which appears to be directed along our line of sight (Schreyer et al. 1997). Shu et al. (1987) hypothesize that (low mass) stars self-determine their final mass through the removal of a collapsing envelope by their outflows. By comparing clusters in different evolutionary states and examining the gas flows around them, it will be possible to see if such a process is relevant in a clustered, massive star forming, environment, and therefore assess its role in the origin of the stellar IMF (Adams & Fatuzzo 1996).

The average HCO\(^+\) spectra toward the IRS2 cluster also shows red-shifted self-absorption and model fits imply a similar infall speed and mass infall rate.
(Williams & Garland 2001). However, the cluster is more widely dispersed and its more irregular appearance precludes a similar radial average study as in IRS1. On the other hand, the greater separation of the protostars allows an investigation of the flows around each one individually. Spectra toward each of six identified cores are shown in Figure 4.

**Figure 4:** Spectra toward individual protostars in IRS2. The central panel shows the continuum emission smoothed to the spectral line resolution of 29″. Contours begin at and are in steps of 0.2 Jy beam⁻¹. Six distinct peaks of emission are labeled a-f and spectra (light: HCO⁺; dark: H¹³CO⁺, multiplied by 3 for clarity) toward each one are shown in the panels on either side.

Core c shows the clearest infall signature with a deep, red-shifted, absorption dip. A model fit using the RATRAN code implies an infall speed $v_{\text{in}} = 0.3 \text{ km s}^{-1}$, and mass infall rate $M_{\text{in}} = 2 \times 10^{-5} \ M_\odot \text{yr}^{-1}$. The speed is the same as the collapse onto the entire cluster but $M$ is a factor of 3 less than the total mass infall rate divided by the number of protostars: apparently not all the material falling onto the cluster makes it all the way to the scale of protostars. This is similar to the case in IRS1 but here we witness the effect on a star-by-star basis as in core d which has a reversed (outflow) spectrum.

The six cores have a dispersion $\sigma_{c-c} = 0.90 \text{ km s}^{-1}$ and extend over an area with an equivalent circular radius $r = 0.44 \text{ pc}$. Together, these imply a virial mass, $M_{\text{virial}} = 3r\sigma_{c-c}^2/G = 250 \ M_\odot$. This is less than the total interstellar mass estimated from the continuum observations but is likely much greater than the final stellar mass of the system (Lada et al. 1993). Thus, although the system is currently bound, as the surrounding gas is dispersed through, e.g., stellar outflows, it is likely to become unbound (Hills 1980).

An unbound cluster expands and its age may be estimated from its size and expansion rate. This effect is seen in optically visible clusters in Orion over a range of ages $\sim 3 - 12 \text{ Myr}$ (Blaauw 1991). The morphology of the two groups in Figure 2 hint that this process may begin even in the protostellar stage: the relatively low column densities toward IRS2 suggest that it is more evolved than...
IRS1 and offer a natural explanation for its greater size. Working backwards from its size and velocity dispersion implies a kinematic age, \( t = \frac{r}{\sigma_{c-c}} = 5 \times 10^5 \, \text{yr} \), characteristic of late Class I sources. Given the continued infall onto the cluster and the large envelope mass, it is not clear if the cluster is actually expanding and whether this simple age estimate is a useful measure at these early stages in a cluster's formation. However, it offers the following simple, testable, picture of cluster evolution.

3. A possible scenario for cluster evolution

Do clusters evolve from a compact state such as IRS1 to a more extended state such as IRS2 as the circumcluster (and circumstellar) envelope is dispersed? Additional observations of other embedded stellar groups in NGC2264 and other regions will show if there is a consistent pattern between cluster age (as judged by the SEDs of the cluster and protostars within) and morphology. In this regard, it is noteworthy that the widely separated group of MSX sources in IRS3 (see Figure 1) was not detected at 870 \( \mu \text{m} \) to the same noise level (0.1 Jy beam\(^{-1}\)) as IRS1 and 2. Thus we might imagine the cartoon model of cluster evolution in Figure 5.

![Figure 5](image)

**Figure 5:** Schematic model of cluster evolution. This simple toy model consists of a group of stars expanding at a constant velocity embedded in individual protostellar envelopes and surrounded by a circumcluster envelope that spreads out linearly with time. The left panel represents a stellar group at early (Class 0) times: protostars are closely spaced and deeply embedded within a common, compact envelope. As the cluster evolves, the envelope begins to disperse and the (Class I) protostars migrate away from each other although there is still substantial circumcluster emission. At later stages, in the rightmost panel, the circumcluster gas has mostly dispersed, and the stars are widely separated, surrounded only by relatively low mass (Class II) protostellar envelopes.

Observations of the early, bright, compact phase of cluster evolution is well suited to interferometry. In later stages, large scale mapping with sub-millimeter bolometer arrays are necessary to image the more extended groups and targeted radio observations to detect remnants of low mass cores around more evolved sources (Saito et al. 2001) can provide essential kinematic information to determine the gravitational stability of a group and estimate its age.
Studies of young clusters are essential for understanding the way in which most, and particularly massive, stars form. The above cartoon model best applies to the accretion model of massive star formation (Maeder; McKee, this volume) but the agglomeration model (Bonnell, this volume) requires an earlier stage, lasting for $\sim 10^5$ yr, consisting of a large group of low mass protostars that first contracts before creating massive stars. Searches for such protoclusters provide a test of the competing theories.

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