Contributions to Polar Amplification in CMIP5 and CMIP6 Models

L. C. Hahn1*, K. C. Armour1,2, M. D. Zelinka3, C. M. Bitz1 and A. Donohoe4

1Department of Atmospheric Sciences, University of Washington, Seattle, WA, United States, 2School of Oceanography, University of Washington, Seattle, WA, United States, 3Lawrence Livermore National Laboratory, Livermore, CA, United States, 4Polar Science Center, Applied Physics Lab, University of Washington, Seattle, WA, United States

As a step towards understanding the fundamental drivers of polar climate change, we evaluate contributions to polar warming and its seasonal and hemispheric asymmetries in Coupled Model Intercomparison Project phase 6 (CMIP6) as compared with CMIP5. CMIP6 models broadly capture the observed pattern of surface- and winter-dominated Arctic warming that has outpaced both tropical and Antarctic warming in recent decades. For both CMIP5 and CMIP6, CO2 quadrupling experiments reveal that the lapse-rate and surface albedo feedbacks contribute most to stronger warming in the Arctic than the tropics or Antarctic. The relative strength of the polar surface albedo feedback in comparison to the lapse-rate feedback is sensitive to the choice of radiative kernel, and the albedo feedback contributes most to intermodel spread in polar warming at both poles. By separately calculating moist and dry atmospheric heat transport, we show that increased poleward moisture transport is another important driver of Arctic amplification and the largest contributor to projected Antarctic warming. Seasonal ocean heat storage and winter-amplified temperature feedbacks contribute most to the winter peak in warming in the Arctic and a weaker winter peak in the Antarctic. In comparison with CMIP5, stronger polar warming in CMIP6 results from a larger surface albedo feedback at both poles, combined with less-negative cloud feedbacks in the Arctic and increased poleward moisture transport in the Antarctic. However, normalizing by the global-mean surface warming yields a similar degree of Arctic amplification and only slightly increased Antarctic amplification in CMIP6 compared to CMIP5.

Keywords: CMIP6, CMIP5, polar amplification, climate feedbacks, Arctic, Antarctic

INTRODUCTION

Observations (Serreze et al., 2009; Screen and Simmonds, 2010a) and climate model projections (Manabe and Stouffer, 1980; Holland and Bitz, 2003) consistently exhibit a pattern of enhanced surface warming in the Arctic compared to the rest of the globe. This so-called Arctic amplification peaks during winter and is at its minimum during summer (Manabe and Stouffer, 1980; Holland and Bitz, 2003; Screen and Simmonds, 2010b; Deser et al., 2010). There is also a strong warming asymmetry between the poles: Antarctic amplification has yet to be observed and is projected to be much weaker than Arctic amplification (Marshall et al., 2015; Smith et al., 2019). Multiple processes contribute to polar amplification, making it a robust feature of the long-term climate response to forcing while at the same time making polar warming inherently more uncertain than global-mean warming (e.g., Holland and Bitz, 2003; Roe et al., 2015; Bonan et al., 2018; Stuecker et al., 2018). Further investigation into the causes of polar
warming and its seasonal and hemispheric asymmetry is thus needed to develop reliable projections of future polar change.

Studies examining a suite of climate models in the Coupled Model Intercomparison Project phase 5 (CMIP5; Block et al., 2020; Goosse et al., 2018; Pithan and Mauritsen, 2014) have quantified key contributors to the magnitude and intermodel spread of polar amplification, motivating the direction of further research with a refined focus. These studies suggest that the largest contributor to Arctic-amplified warming is the lapse-rate feedback, which is more positive in the Arctic than elsewhere. Unlike in the tropics, where deep convection causes surface warming to be amplified with height, in polar regions, a stable lower troposphere inhibits vertical mixing and contributes to stronger warming near the surface than aloft (Cronin and Janssen, 2015; Payne et al., 2015; Hahn et al., 2020). This surface-trapped warming leads to a positive lapse-rate feedback by producing less longwave emission to space than a vertically uniform heating of the atmospheric column.

While the positive surface albedo feedback associated with sea-ice loss plays a key role in polar warming (e.g., Hall, 2004; Dai et al., 2019), Pithan and Mauritsen (2014) suggest that its contribution to Arctic amplification is secondary to that of the lapse-rate feedback. However, the strength of the lapse-rate feedback itself is highly dependent on the degree of surface warming and sea-ice loss (Graversen et al., 2014; Fedl et al., 2017; Fedl et al., 2020; Boeke et al., 2021). The albedo feedback additionally contributes most to intermodel spread in polar warming among CMIP5 models, followed by the lapse-rate feedback (Pithan and Mauritsen, 2014).

Another substantial contribution to Arctic amplification is made by the Planck feedback (Pithan and Mauritsen, 2014). Following the Stefan-Boltzmann law, a given surface warming at initially colder temperatures produces a weaker increase in emitted longwave radiation, causing a less-negative Planck feedback in the Arctic than the tropics. Finally, poleward atmospheric heat transport (AHT) into the Arctic increases only a small amount under climate warming within CMIP5 models, suggesting that AHT makes only a small contribution to Arctic amplification (Pithan and Mauritsen, 2014; Goosse et al., 2018). However, many studies highlight increased latent heat transport into the Arctic as a primary driver of polar warming and note that the small change in total AHT reflects compensating changes in latent and dry heat transports (e.g., Alexeev et al., 2005; Armour et al., 2019). Other processes such as water-vapor and cloud feedbacks, Arctic surface heat fluxes (i.e., ocean heat uptake), and the meridional structure of CO₂ forcing contribute more to tropical than polar warming (Pithan and Mauritsen, 2014; Goosse et al., 2018).

Assessments of polar warming in CMIP5 highlight key drivers not only of Arctic amplification, but also of seasonal and hemispheric asymmetry in polar warming. Summer ocean heat storage and its release to the atmosphere in winter contributes most to the winter peak in Arctic warming, and drives substantial intermodel spread in Arctic amplification (Pithan and Mauritsen, 2014; Boeke and Taylor, 2018). As a result of enhanced vertical stability in winter, the winter-peaking lapse-rate feedback makes an additional contribution to the winter maximum in Arctic warming (Pithan and Mauritsen, 2014). The lapse-rate feedback is also the largest contributor to greater warming in the Arctic than Antarctic in CMIP5 models (Goosse et al., 2018) due to the elevation of the Antarctic ice sheet and resulting shallower and weaker base-state Antarctic inversions (Salzmann, 2017; Hahn et al., 2020). Goosse et al. (2018) also confirm a large role for Southern Ocean heat uptake (Marshall et al., 2015; Armour et al., 2016) and a more-negative cloud feedback in weakening transient Antarctic warming compared to the Arctic.

The Coupled Model Intercomparison Project phase 6 (CMIP6) offers an opportunity to reexamine the processes contributing to polar amplification in a new model ensemble and evaluate the evolution of relevant processes between model generations. The higher effective climate sensitivity in CMIP6 than CMIP5, on average, has been traced to less-negative extratropical cloud feedbacks within many CMIP6 models (Zelinka et al., 2020), suggesting that extratropical cloud feedbacks may also contribute more to polar warming in CMIP6. To explore how cloud and other feedbacks contribute to polar amplification in CMIP6 models, we apply a ‘warming contribution’ analysis (Pithan and Mauritsen, 2014; Goosse et al., 2018) to CMIP6 and compare with the same analysis applied to CMIP5. We evaluate the drivers of Arctic amplification, weaker Antarctic amplification, and seasonal asymmetry in polar warming, considering also the spread in warming contributed by model differences. We note that Cai et al. (2021) have previously examined Arctic and Antarctic warming contributions in CMIP6, but here we expand this analysis from 15 to 42 CMIP6 models (Supplementary Table S1), add a comparison to CMIP5, consider climate feedback sensitivity to the choice of radiative kernel, and consider more closely the role of AHT in driving polar amplification by partitioning its changes into moist and dry components. We also apply the warming contribution analysis to Atmospheric Model Intercomparison Project phase 6 (AMIP6) models to estimate contributions to historical modelled warming in comparison to warming projected by fully-coupled CMIP6 models. By quantifying key contributors to polar warming and its asymmetries in CMIP6, we hope to assess previously established mechanisms of Arctic amplification as well as identify open questions in support of future polar research.

**HISTORICAL POLAR AMPLIFICATION IN OBSERVATIONS AND MODELS**

Before quantifying contributions to projected surface warming in CMIP6, we compare historical near-surface and atmospheric temperature trends over 1979–2014 from fully-coupled CMIP6 models with reanalysis data and observations (Figures 1, 2). We use the European Centre for Medium-Range Weather Forecasts Interim Re-Analysis (ERA-Interim; Dee et al., 2011) which, in an evaluation of seven reanalyses over the Arctic, has been found to perform best in simulating observations of near-surface air temperature, surface radiative fluxes, precipitation, and wind speed (Lindsay et al., 2014). Limitations of this and other reanalyses in the Arctic include a positive near-surface
temperature bias over sea ice in winter, with a slightly smaller bias for ERA-Interim than its successor, ERA5 (Graham et al., 2019; Wang et al., 2019). Although ERA-Interim also overestimates the lowest observed near-surface temperatures over Antarctica, it correlates relatively well with Antarctic observations of near-surface temperatures (Gossart et al., 2019) and their trends

**FIGURE 1** Annual-mean near-surface temperature trends (°C decade$^{-1}$) for 1979–2014 from (A) HadCRUT5 observations, (B) the ERA-Interim reanalysis, and (C) the historical CMIP6 multimodel mean.
We also show historical near-surface temperature trends for HadCRUT5, an observational dataset which combines CRUTEM5 near-surface air temperature over land with HadSST4 sea-surface temperatures, with statistical infilling where observations are unavailable (Morice et al., 2021). Throughout this study, we define the Arctic as 60 to 90°N and the Antarctic as 60 to 90°S, and define polar amplification as the near-surface warming poleward of 60° in a given hemisphere divided by global-mean near-surface warming.

The ensemble mean of fully-coupled CMIP6 models (hereafter called the CMIP6 mean) for historical simulations reproduces the observed pattern of amplified Arctic warming and weaker warming in the Antarctic, but exceeds the observed warming at all latitudes except in the Arctic (Figures 1, 2A). As a result, the CMIP6 mean produces too little Arctic amplification and too much Antarctic amplification over this historical period: the degree of Arctic and Antarctic amplification in the CMIP6 mean is 2.6 and 0.9, respectively, compared to 3.5 and 0.4 in HadCRUT5 observations. Figure 2 also shows the CMIP6 mean near-surface temperature trend with ± 2 standard deviations of trends calculated across the first 90 ensemble members of the Community Earth System Model version 2 Large Ensemble (CESM2-LE) as a reference for the range of internal variability within a single climate model in CMIP6. The CESM2-LE uses identical time-varying external forcing for its ensemble members, but begins each with a combination of different oceanic and atmospheric initial conditions (Rodgers et al., 2021). With the exception of the Southern Hemisphere mid-latitudes, zonal-mean warming trends in HadCRUT5 fall within the range of CMIP6 intermodel spread and the range of internal variability for CESM2-LE, suggesting that CMIP6 models may be largely consistent with observations when internal variability is taken into account.

CMIP6 models capture the seasonality of near-surface warming in the Arctic, with a peak in warming during early winter (Figure 2B). While the CMIP6 multimodel mean excludes a second peak in warming in April found in observations and reanalyses (Screen et al., 2012), this April maximum falls within the intermodel spread for CMIP6 and the ensemble spread for CESM2-LE, suggesting that it may be explained by internal variability. Model underestimation of observed reductions in spring snow cover extent over land may also contribute to a low bias in modeled Arctic springtime warming (Screen et al., 2012; Brutel-Vuilmet et al., 2013). In the Antarctic, CMIP6 models on average simulate year-round warming with a winter maximum (Figure 2C) whereas the observations and reanalysis show near-zero warming or slight cooling. Both the intermodel spread and spread due to internal variability in historical near-surface warming are largest in polar regions and during winter.

In addition to considering historical trends from fully coupled models in CMIP6, we show historical near-surface and atmospheric temperature trends from models in the
FIGURE 3 | Atmospheric temperature trends (°C decade\(^{-1}\)) for 1979–2014 for (A,C,E) the Arctic and (B,D,F) the Antarctic in the ERA-Interim reanalysis (A,B), and historical CMIP6 (C,D) and AMIP6 (E,F) multimodel mean trends. Black dots in (A,B) show statistically significant trends at the 95% level based on a two-tailed Student’s t-test, and black dots in (C,D) and (E,F) show where 75% of models meet these criteria for significant trends. For the reanalysis and models, temperature data is first masked at pressures greater than the surface pressure, and trends are only shown where more than 50% of area-averaged grid points are non-missing.
Atmospheric Model Intercomparison Project phase 6 (AMIP6; Supplementary Table S1). These models prescribe time-varying monthly sea-surface temperatures and sea-ice concentrations based on observations, while still including atmosphere-land coupling, along with time-varying historical forcing agents (Eyring et al., 2016). The ensemble mean of AMIP6 models aligns more closely than the CMIP6 mean with observed near-surface atmospheric temperature trends in the tropics and Antarctic (Figures 2D–F). Despite having prescribed observational sea-surface temperature and sea-ice concentration changes, the AMIP6 mean underestimates observed Arctic warming over both land and sea-ice surfaces. This may result from differences between models and observations in sea-ice thickness and the near-surface air temperature response to sea-ice loss. Underestimated Arctic amplification in both the CMIP6 and AMIP6 means is also shown in Supplementary Figure S1, with historical temperature trends normalized by global-mean near-surface warming.

Consistent with previous modelling and observational studies, the vertical and seasonal pattern of Arctic warming for both ERA-Interim and CMIP6 is amplified near the surface during winter and is more vertically-uniform during summer (Figures 3A,C). For both summer and winter, mid-tropospheric warming trends are stronger for CMIP6 than ERA-Interim, which may reflect stronger increased poleward AHT due to overestimated mid-latitude surface warming (Fajber et al., 2018; Feldl et al., 2020; Kosaka and Xie, 2013). Consistent with this hypothesis, AMIP6 models, which simulate mid-latitude near-surface temperature trends closer to observations, show weaker warming aloft than CMIP6 (Figure 3E). AMIP6 models still demonstrate stronger mid-tropospheric summer warming than ERA-Interim, which may be related to AHT differences contributed by surface temperatures over land, or to enhanced atmospheric shortwave absorption by water vapor in these models (Donohoe and Battisti, 2013).

In contrast to CMIP6 models, which propagate Arctic winter surface warming further aloft than the ERA-Interim reanalysis, CMIP3 and CMIP5 models simulate excessively surface-trapped Arctic warming as a result of overestimating mean-state inversion strength (Medeiros et al., 2011; Pithan et al., 2014; Screen et al., 2012). This overestimated stability in earlier models has been attributed to a low bias in the supercooled liquid fraction of mixed-phase clouds, which allows for too much surface radiative cooling (Pithan et al., 2014). While CMIP6 models also tend to overestimate winter surface inversion strength in the Arctic (Supplementary Figure S2), increased supercooled liquid fraction in some CMIP6 models may have reduced these biases in inversion strength and surface-trapped warming compared to earlier models. A weaker ice-to-liquid transition in CMIP6 models with more mean-state supercooled liquid may also lead to a smaller increase in downward longwave radiation under warming, additionally causing weaker surface-trapped warming in CMIP6 (Tan and Storelvmo., 2019).

In both models and observations, Antarctic temperature trends are much weaker than those in the Arctic (Figures 3B,D,F; note the reduced colorbar range for the Antarctic). CMIP6 models demonstrate a small surface-amplified warming during winter, while ERA-Interim shows surface-amplified cooling particularly for December-May, as observed in near-surface temperature trends. In summary, although differences exist between observed and modeled Antarctic temperature trends, CMIP6 models generally agree with observations in producing winter- and surface-amplified warming in the Arctic, with much weaker trends in the Antarctic. Next we investigate the drivers of this Arctic-amplified warming and its hemispheric and seasonal asymmetry in CMIP5 and CMIP6 models.

**CONTRIBUTIONS TO POLAR WARMING IN CMIP5 AND CMIP6 MODELS**

**Warming Contribution Methodology**
Following previous studies (Lu and Cai, 2009; Crook and Forster 2011; Feldl and Roe 2013; Taylor et al., 2013; Pithan and Mauritsen, 2014; Goosse et al., 2018), we calculate contributions to projected polar warming based on an energy budget analysis. To do so, we use CMIP5 and CMIP6 output from pre-industrial control (piControl) and abrupt CO2 quadrupling (abrupt4xCO2) experiments, in which CO2 concentrations are quadrupled from piControl conditions and then held fixed for 150 years. As in Zelinka et al. (2020) and Caldwell et al. (2014), we apply a 21-year running average to piControl experiments to account for model drift before computing anomalies between abrupt4xCO2 and piControl during corresponding time periods (i.e., after branching to abrupt4xCO2, year-100 in piControl is compared to year-100 in abrupt4xCO2). With the assumption of similar model drift in the abrupt4xCO2 and piControl experiments, this approach isolates anomalies due to CO2 forcing rather than model drift. We calculate the effective radiative forcing (ERF) as the y-intercept of the regression between top-of-atmosphere (TOA) radiation anomalies at each grid point against the global-mean near-surface temperature anomalies for the first 20 years after CO2 quadrupling (Gregory et al., 2004). Smith et al. (2020) demonstrate that this 20-year regression yields ERF values which closely match methods using fixed sea-surface temperatures (Hansen et al., 2005) under CO2 quadrupling in CMIP6 models, while regression over the full 150-year abrupt4xCO2 period instead underestimates ERF as a result of time-varying feedbacks in models (Andrews et al., 2015; Dong et al., 2020).

To calculate temperature anomalies, climate feedbacks, and heat transport anomalies under CO2 quadrupling, we use monthly climate variable anomalies averaged over 31 years centered on year-100 of the abrupt4xCO2 experiments. We calculate climate feedbacks using the radiative kernel method (Shell et al., 2008; Soden et al., 2008), in which relevant climate variable anomalies are multiplied by monthly- and spatially-resolved radiative kernels, which quantify the change in radiative flux per unit change in a given climate variable. Vertically integrating this product throughout the troposphere gives the contribution of each feedback to TOA radiation anomalies. Temperature feedbacks are separated into the effect
of surface temperature changes propagated throughout the troposphere (the Planck feedback) and the effect of departures from this vertically uniform temperature change (the lapse-rate feedback). Cloud feedbacks are calculated using the change in cloud radiative forcing (ΔCRF), equal to the change in all-sky minus clear-sky TOA radiation, minus a cloud masking term. This cloud masking term is defined as the effect of noncloud variables (temperature, water vapor, and surface albedo) on ΔCRF, calculated using all-sky and clear-sky radiative kernels. While our method is consistent with Goosse et al. (2018), one caveat of using year-100 feedback energetic contributions divided by year-100 temperature anomalies as opposed to linear regression of these fields (as in Zelinka et al. (2020)) is that the resulting feedbacks include both the true temperature-mediated feedbacks and rapid adjustments that occur immediately upon quadrupling CO₂. Resulting differences from the regression method used in Zelinka et al. (2020) for cloud feedbacks are discussed at the end of the section Warming Contributions in CMIP5 and CMIP6.

We primarily show feedbacks calculated using the Huang et al. (2017) kernels, which are based on ERA-Interim reanalysis data and give the smallest residual terms between the modelled TOA radiation anomalies and the sum of feedback and forcing radiation contributions (Zelinka et al., 2020). However, we additionally consider sensitivity to kernel choice, including kernels from Soden et al. (2008) (based on GFDL AM2), Shell et al. (2008) (based on NCAR CAM3), Block and Mauritsen (2013) (based on NCAR CAM5), Smith et al. (2018) (based on HadGEM2). For comparison with the kernel-derived surface albedo feedback, we also compute the surface albedo feedback using the approximate partial radiative perturbation (APRP) method (Taylor et al., 2007).

As in Pithan and Mauritsen (2014), we calculate the annual AHT convergence as the difference between surface and net TOA fluxes. We further partition this into a moist component using the difference between precipitation and evaporation multiplied by the latent heat of vaporization (with the latent heat of fusion for solid precipitation neglected), and a dry component calculated as the residual between total and moist AHT convergence. To calculate the seasonal cycle of AHT convergence, we additionally subtract atmospheric energy and moisture storage terms following Donohoe et al. (2020a). Anomalous surface heat fluxes (referred to here as ocean heat uptake) implicitly include both ocean heat transport and ocean heat storage, on both seasonal and annual timescales.

We use a local energy budget (Eq. 1 below) to convert these energetic contributions of climate feedbacks and heat transport anomalies surrounding year-100 of the abrupt4xCO2 experiments into contributions to near-surface warming (ΔT) in the tropics, Arctic, and Antarctic, as in previous studies (Lu and Cai, 2009; Crook and Forster, 2011; Feldl and Roe, 2013; Taylor et al., 2013; Pithan and Mauritsen, 2014; Goosse et al., 2018). Equation 1 includes the ERF, energetic contributions of climate feedbacks (λpΔT) and the Planck response (λpΔT), anomalies in AHT convergence (AAHT) and ocean heat uptake (ΔO), and a residual term (ΔRres), all in units of W m⁻²:

\[ ERF + \left( \sum_{i} \lambda_i \right) \Delta T + \Delta AHT + \Delta O + \Delta R_{res} = 0 \] (1)

In addition to computing annual-mean warming contributions, we calculate contributions during winter (December-January-February for the Arctic and June-July-August for the Antarctic) and summer (June-July-August for the Arctic and December-January-February for the Antarctic) seasons. For each region and season, warming contributions are defined by dividing each term in Eq. 1 by the global- and annual-mean Planck feedback (λp) in W m⁻² K⁻¹:

\[ \Delta T = -\frac{ERF}{\lambda_p} - \frac{\lambda_p \Delta T}{\lambda_p} - \sum_i \frac{\lambda_i \Delta T}{\lambda_p} - \frac{\Delta AHT}{\lambda_p} - \frac{\Delta O}{\lambda_p} - \frac{\Delta R_{res}}{\lambda_p} \] (2)

where \( \lambda_p = \lambda_p - \bar{\lambda_p} \) is the difference between the regional, seasonal Planck feedback, \( \lambda_p \), and its annual- and global-mean value, \( \bar{\lambda_p} \).

### Warming Contributions in CMIP5 and CMIP6

Near-surface temperature anomalies centered around year-100 of the abrupt4xCO2 simulations in CMIP5 and CMIP6 are shown in Figure 4A. Consistent with observed and modelled historical temperature trends (Figures 1, 2), both CMIP5 and CMIP6 models project transient warming under CO₂ quadrupling that is amplified in the Arctic compared to the tropics, with weaker Antarctic amplification. CMIP6 models exhibit large intermodel spread in the Arctic, with an interquartile range of up to 8°C in Arctic warming compared to about 4°C in the Antarctic and 2°C in the tropics. In the multimodel mean, Arctic warming has increased from 10.1°C in CMIP5 to 11.5°C in CMIP6, while Antarctic warming has increased from 5.1°C in CMIP5 to 6.4°C in CMIP6. This increase in polar warming is greater than the increase in tropical warming from 4.3°C in CMIP5 to 4.8°C in CMIP6. However, normalizing by the global-mean surface warming in each model (Figure 4B) demonstrates a similar degree of Arctic amplification and only slightly increased Antarctic amplification in CMIP6 compared to CMIP5.

To investigate the drivers of polar amplification and hemispheric asymmetry in CMIP5 as compared to CMIP6, we calculate contributions to polar warming from feedbacks, AHT changes, and ocean heat uptake following Eq. 2. We use the Huang et al. (2017) kernels for climate feedbacks in Figures 4C,D, and investigate the sensitivity to kernel choice in Figure 5. Consistent with Pithan and Mauritsen (2014) and Goosse et al. (2018), key contributors to Arctic amplification in both CMIP5 and CMIP6 are the lapse-rate, albedo, and Planck feedbacks (Figure 4C). In contrast to the secondary role of the albedo feedback found in Pithan and Mauritsen (2014) using the Block and Mauritsen (2013) radiative kernels, use of the Huang et al. (2017) kernels yields lapse-rate and albedo feedbacks of almost equal importance for polar amplification in CMIP5, and equivalent importance in CMIP6. Partitioning AHT into moist and dry components illustrates that while reduced dry AHT opposes Arctic amplification, increased moist AHT is a large contributor to Arctic amplification.
Stronger Arctic warming in CMIP6 than CMIP5 is mainly contributed by more-positive albedo and less-negative cloud feedbacks. Less-negative Arctic cloud feedbacks in CMIP6 result from less-negative shortwave low cloud amount and scattering feedbacks, likely due to updated treatment of supercooled liquid fraction in mixed phase clouds (Zelinka et al., 2020). The lapse-rate feedback, Planck response, and moist AHT changes also contribute to stronger Arctic warming in CMIP6, while increased ocean heat uptake and equatorward dry AHT more strongly oppose Arctic warming in CMIP6. Normalizing Arctic warming contributions by the global-mean warming yields contributions which sum to the total Arctic amplification in CMIP5 and CMIP6 (Supplementary Figure S3A, vertical axis): less-negative cloud and more-positive albedo feedbacks support greater Arctic amplification in CMIP6, while most other contributions support weaker Arctic amplification in CMIP6 due to normalizing by the larger global-mean warming in CMIP6 than CMIP5. This results in a similar degree of Arctic amplification in CMIP5 and CMIP6 (Figure 4B).

Consistent with Goosse et al. (2018), the largest contributor to stronger warming in the Arctic than Antarctic is the lapse-rate feedback for both CMIP5 and CMIP6 (Figure 4D). In fact, all factors except for CO2 forcing and moist AHT changes support greater warming in the Arctic than Antarctic, with an additionally large contribution from the albedo feedback in CMIP5 and CMIP6. This feedback asymmetry between the poles is supported by the elevation of the Antarctic ice sheet (Salzmann, 2017), which primarily weakens the Antarctic lapse-rate feedback through reducing the average strength of mean-state inversions (Hahn et al., 2020). We note that the Planck feedback (in W m$^{-2}$ K$^{-1}$) is slightly less negative in the Antarctic than Arctic, likely due to colder and drier initial conditions, but that the Planck warming contribution is larger in the Arctic due to a larger Arctic $\Delta T$ resulting in a larger contribution $\lambda' \Delta T$ in Eq. 2. This illustrates one limitation of the warming contribution framework: a warming contribution from one feedback is influenced by all other feedbacks through their influence on $\Delta T$. 

**FIGURE 4 |** (A) Annual- and zonal-mean near-surface warming (˚C) averaged over 31 years centered on year-100 after CO2 quadrupling for the CMIP6 (solid orange line) and CMIP5 (dashed orange line) multimodel means. The dark orange shading shows the 25th to 75th percentiles, and the light orange shading shows the full intermodel spread for CMIP6. (B) As in (A), but with zonal-mean near-surface warming normalized by global-mean near-surface warming within each model. (C,D) Contributions of each feedback and atmospheric forcing to warming (˚C) centered around year-100 of abrupt CO2 quadrupling in CMIP6 (filled circles) and CMIP5 (hollow circles) for (C) the tropics relative to the Arctic and (D) the Arctic relative to the Antarctic. Warming contributions are shown for the lapse-rate (LR), surface albedo (A), water-vapor (WV), and cloud (C) feedbacks, the variation in the Planck response from its global-mean value ($P'$), effective radiative forcing ($CO_2$), change in moist and dry AHT convergence ($\Delta AHT_m, \Delta AHT_d$) and ocean heat uptake (Ocean), and residual term (Res). Dashed grey line shows a 1-to-1 slope through the lapse-rate feedback warming contribution.
Moist AHT is the largest contributor to Antarctic warming in both CMIP5 and CMIP6. Combined with initially colder Antarctic temperatures and Clausius-Clapeyron nonlinearity, weaker warming in the Antarctic under CO₂ quadrupling produces a weaker moisture increase compared to the Arctic. As a result, the equator-to-Antarctic latent heat gradient increases more than the equator-to-Arctic gradient, contributing to a stronger increase in moist AHT to the Antarctic (Supplementary Figure S4). Moist AHT changes are also sensitive to climate feedbacks which alter the equator-to-pole moist static energy gradient, particularly shortwave cloud feedbacks (Hwang and Frierson, 2010; Zelinka and Hartmann, 2012; Shaw and Voigt, 2016; Chen et al., 2021). More-negative shortwave cloud feedbacks in the Antarctic may therefore also contribute to larger increased moist AHT to the Antarctic than Arctic by enhancing the equator-to-pole moist static energy gradient in the Southern Hemisphere.

Most warming contributions change similarly for the Arctic and Antarctic from CMIP5 to CMIP6 (i.e., on a one-to-one slope in Figure 4D), including the lapse-rate, water-vapor, Planck, and albedo feedbacks, as well as ocean heat uptake and dry AHT. The two warming contributions that change differently for the Arctic and Antarctic from CMIP5 to CMIP6 are the cloud and moist AHT contributions. While the cloud contribution primarily increases in the Arctic, the moist AHT contribution primarily increases in the Antarctic. Stronger Arctic than Antarctic changes in cloud feedbacks result from shortwave cloud feedback changes (Supplementary Figure S5), and this polar difference appears to be amplified by the use of year-100 feedbacks rather than the 150-year regression method of Zelinka et al. (2020). Still, the year-100 feedbacks are generally consistent with the 150-year regression method in demonstrating less-negative polar cloud feedbacks in CMIP6 compared to CMIP5 (Supplementary Figure S5).

**Dependence on Choice of Kernel and Feedback Definition**

While most feedbacks are relatively insensitive to the choice of radiative kernel, polar surface albedo and cloud feedbacks particularly for the Arctic show substantial kernel sensitivity in both CMIP5 and CMIP6 (Figure 5). This is consistent with evidence that the radiative sensitivity to albedo changes (the albedo radiative kernel) varies by a factor of two across climate
models in the Arctic and Southern Ocean due to intermodel differences in mean-state cloudiness (Donohoe et al., 2020b). Kernel sensitivity in the albedo feedback also contributes to kernel sensitivity in the cloud feedback, which is calculated using radiative kernels to compute and subtract the cloud masking effect of noncloud variables, including surface albedo, from the total ΔCRF. The APRP method gives an albedo feedback near the bottom of the range in kernel-derived albedo feedbacks. This may result from using the average of a forward- and backward-radiative substitution in the APRP method, whereas the kernels rely solely on a forward calculation. The surface albedo feedback derived from the APRP method versus radiative kernels is thus conceptually different quantities, as the APRP method allows cloud changes to impact the surface albedo feedback while the kernel method does not.

Of the model-derived surface albedo kernels, the Smith et al. (2018) kernels come closest to simulating the radiative sensitivity to albedo changes derived from satellite observations in the Arctic (Donohoe et al., 2020b). The Smith et al. (2018) kernels also produce an Arctic albedo feedback similar to the Huang et al. (2017) observationally-derived kernels (Figure 5A). This suggests that the observed mean state is consistent with a stronger Arctic albedo feedback than previously found, on par with the lapse rate feedback in its contribution to Arctic amplification.

An important result of kernel sensitivity in the albedo feedback is that the relative importance of the albedo versus lapse-rate feedback depends on the choice of kernel. However, for all kernels the lapse-rate and albedo feedbacks remain key contributors to Arctic amplification and hemispheric asymmetry in polar warming. Additionally, because the albedo and cloud feedbacks have compensating sensitivity to kernel choice, the total polar feedback remains relatively insensitive to kernel choice, as evidenced by the small kernel sensitivity in the residual term.

In addition to the traditional feedback framework applied here, alternative feedback definitions can be used, including a framework which quantifies the effect of warming and moistening at constant relative humidity (RH) separately from the effect of RH changes (Held and Shell, 2012). We compare the traditional feedback framework with the fixed-RH method, where the lapse-rate and Planck feedbacks are calculated at constant RH, an RH feedback is calculated, and all other feedbacks are identical to the traditional feedbacks (Supplementary Figure S6). Consistent with Held and Shell (2012), the magnitude of the fixed-RH Planck, fixed-RH lapse-rate, and RH feedbacks is reduced compared to the traditional Planck, lapse-rate, and water-vapor feedbacks. Although the other feedbacks are unchanged, division by a weaker global Planck feedback contributes to larger fixed-RH warming contributions for these feedbacks. Applied to the tropics, the fixed-RH framework gives a less-negative lapse-rate feedback than in the traditional framework: while amplified warming aloft promotes a large, negative lapse-rate feedback, amplified moistening aloft to maintain constant relative humidity offsets this negative feedback. In the Arctic, the fixed-RH framework produces a less-positive lapse-rate feedback: while weaker warming aloft compared to the surface supports a large, positive lapse-rate feedback, weaker moistening aloft to maintain constant relative humidity reduces the magnitude of the positive fixed-RH lapse-rate feedback. As a result, the relative contribution of the lapse-rate feedback to Arctic amplification is weakened in the fixed-RH framework, with stronger contributions from the albedo feedback and poleward moisture transport. While feedback definition choice can impact the relative roles of contributions to Arctic warming, we note that moist AHT and the albedo and lapse-rate feedbacks remain important contributors to Arctic amplification for both the traditional and fixed-RH frameworks.

**Intermodel Spread**

Following Pithan and Mauritsen (2014), we also investigate what factors contribute to substantial intermodel spread in polar warming by analyzing intermodel spread in CMIP6 warming contributions in both the Arctic and Antarctic (Figure 6). The albedo feedback is the single largest contributor to increased intermodel spread in both Arctic and Antarctic warming. While the dry AHT term itself exhibits substantial intermodel spread, it contributes more cooling to models with stronger polar warming and thus reduces intermodel spread in polar warming, as shown in Hwang et al. (2011). In contrast, changes in moist AHT generally increase with total polar warming and contribute to intermodel spread. Relationships between total polar warming and each warming contribution are similar for CMIP6 and CMIP5 (Supplementary Figure S7) with the exception of ocean heat uptake changes. In CMIP5, the ocean term becomes more negative (greater ocean heat uptake) in models with greater Arctic warming, while in CMIP6, models with weaker ocean heat uptake simulate greater Arctic warming. In the Antarctic, CMIP5 models with weaker ocean heat uptake simulate greater warming, while there is no correlation between ocean heat uptake and Antarctic warming across different models in CMIP6.

Intermodel spread in polar warming is contributed not only by individual warming contributions, but also by their covariances; to quantify both, we show covariance matrices of the contributions to Arctic and Antarctic warming in Figure 7, following Caldwell et al. (2014). Each term has been normalized by the total warming variance (10.2 K² in the Arctic; 4.9 K² in the Antarctic) to illustrate fractional contributions to warming variance in each region. To avoid showing redundant information in these symmetric matrices, covariance terms below the diagonal are omitted and those above the diagonal are multiplied by two. Consistent with Figure 6, the main diagonal in Figure 7 shows large variances contributed by the albedo feedback and dry AHT at both poles. However, strong negative covariance between these two terms leads to a large damping of intermodel spread. Negative covariances between dry AHT and almost every other warming contribution suggest that dry AHT responds to polar warming, with stronger polar warming weakening the equator-to-pole temperature gradient and reducing dry AHT to polar regions. In contrast, large positive covariance between the albedo and lapse-rate feedbacks magnifies the intermodel spread in polar warming. In the Antarctic, variance in moist AHT and its covariance with the albedo
feedback also contribute strongly to total warming variance. Although the total warming variance is smaller in the Antarctic than the Arctic, the albedo feedback constitutes a larger fraction of the total variance in Antarctic warming. These results support previous suggestions that constraining the albedo feedback may reduce intermodel spread in polar warming contributed both directly by this feedback and by covariances with other feedbacks (e.g., Feldl et al., 2020; Boeke et al., 2021).

Seasonality in Polar Warming Contributions
Lastly, we consider what drives seasonality in warming for the Arctic and Antarctic. As seen in historical CMIP6 trends, polar warming under CO$_2$ quadrupling peaks during winter (Figures 8A,B). Compared to the Antarctic, stronger seasonality in Arctic warming largely stems from stronger winter warming, while summer warming is more similar between the poles. Contributions to Arctic seasonality in warming in CMIP6 are consistent with CMIP5 results (Supplementary Figure S8 and Pithan and Mauritsen, 2014): while the albedo and water-vapor feedbacks support stronger summer warming, summer ocean heat storage and its release to the atmosphere in winter contributes to stronger winter warming. In addition, the lapse-rate and Planck feedbacks contribute to winter-amplified Arctic warming. While similar factors contribute to Antarctic seasonality in warming, weaker winter warming in the Antarctic compared to the Arctic results from weaker temperature feedbacks and seasonal ocean heat storage.

Comparison With Historical Warming Contributions in AMIP6
Applying the above methodology to historical AMIP6 simulations allows us to evaluate polar warming contributions within models that use the observed patterns of sea-surface temperatures and sea-ice concentrations as boundary
FIGURE 7 | Fractional contributions of each warming contribution term to intermodel variance in (A) Arctic and (B) Antarctic warming in CMIP6.

FIGURE 8 | Monthly near-surface warming (°C) centered around year-100 of abrupt CO₂ quadrupling for the CMIP6 multimodel mean (orange line), the 25th to 75th percentile (dark orange shading), and the full intermodel spread (light orange shading) in (A) the Arctic and (B) the Antarctic. (C,D) Contributions to winter and summer warming (°C) centered around year-100 of abrupt CO₂ quadrupling in CMIP6 for (C) the Arctic and (D) the Antarctic. Warming contributions are shown for the lapse-rate (LR), surface albedo (A), water-vapor (WV), and cloud (C) feedbacks, the variation in the Planck response from its global-mean value (P), effective radiative forcing (CO₂), change in moist and dry AHT convergence (ΔAHT_m, ΔAHT_d) and ocean heat uptake (Ocean), and residual term (Res).
conditions, which may produce some differences from the fully-coupled CMIP6 results under CO₂ quadrupling shown above. To calculate feedbacks in AMIP6, we compute monthly anomalies in climate variables with respect to the 1979–2014 climatology, and regress radiative contributions of feedbacks against near-surface air temperature anomalies for this period. We then calculate warming contributions again using Eq. 2, where Δ now indicates the trend in each variable from 1979 to 2014, multiplied by the period of 36 years.

Unlike the idealized CO₂ quadrupling experiments, AMIP simulations have time-evolving effective radiative forcing (ERF) that must be accounted for. Previous studies have derived the historical ERF in AMIP6 models using experiments from the Radiative Forcing Model Intercomparison Project (RFMIP, Pincus et al., 2016) with time-varying forcing applied on top of constant pre-industrial sea-surface temperature and sea-ice concentrations (e.g., Zhang et al., 2020). However, these RFMIP experiments are only available for 7 CMIP6 models. To increase our model sample size, we estimate ERF using kernels in each model as follows. Because clear-sky TOA radiation anomalies are equal to the sum of clear-sky feedback energetic contributions, the clear-sky ERF, and a residual term arising from errors in the kernel approach, we estimate the clear-sky ERF as the difference between TOA radiation anomalies and the sum of kernel-derived clear-sky feedback energetic contributions. The neglect of kernel residual terms is justifiable by the fact that 1) kernel-derived and RFMIP-derived estimates of clear-sky ERF are in excellent agreement and 2) kernel residuals are very close to zero in amip-piForcing experiments from the Cloud Feedback Model Intercomparison Project (CFMIP, Webb et al., 2017) in which forcings are held constant at pre-industrial levels while sea-surface temperature and sea-ice concentration fields are prescribed to follow time-varying observations (not shown). Following standard practice, we then estimate the all-sky ERF by dividing clear-sky ERF by 1.16 (Soden et al., 2008). This method allows us to include 38 AMIP6 models in this analysis (Supplementary Table S1).

As in CMIP6 CO₂-quadrupling experiments, historical AMIP6 experiments show strong contributions to Arctic amplification from increased moist AHT and the lapse-rate, Planck, and albedo feedbacks, while the water-vapor feedback, Arctic ocean heat uptake, longwave forcing, and changes in dry AHT oppose Arctic amplification (Figure 9A). Despite differences in the vertical structure of warming between CMIP6 and AMIP6 models, their lapse-rate contributions to Arctic amplification appear relatively similar and consistently on par with respective surface albedo contributions. As suggested by Boeke et al. (2021), this may indicate the strong dependence of the lapse-rate feedback on the surface albedo feedback and surface temperature changes, more so than the vertical structure of warming. Differences in warming contributions between AMIP6 and CMIP6 include a relatively larger Planck contribution to Arctic amplification in AMIP6 and an Arctic-amplified SW ERF contribution in AMIP6. This positive SW forcing may be driven by reduced European sulfate emissions since 1980, which disproportionately warmed the Arctic compared to the rest of the globe (Acosta Navarro et al., 2016).

In both AMIP6 and CMIP6, the lapse rate, water vapor, Planck, and albedo feedbacks contribute to weaker warming in the Antarctic than Arctic, while increased poleward moisture transport contributes more strongly to Antarctic warming (Figure 9B). In contrast to CMIP6 projections, negative Antarctic warming contributions in AMIP6 for the albedo feedback and Planck feedback deviation from its global-mean value reflect historical cooling and sea-ice expansion over the Southern Ocean. As a result, the albedo feedback contributes most to stronger Arctic than Antarctic warming in AMIP6, while the lapse-rate feedback makes the largest contribution to this hemispheric asymmetry in CMIP6 projections. Weaker historical than projected Antarctic warming also weakens the equatorward dry AHT opposing Antarctic warming in AMIP6. These differences between AMIP6 and CMIP6 illustrate the strong
dependence of Antarctic feedbacks on changes in Southern Ocean sea-surface temperature and sea ice.

Even with identical prescribed sea-surface temperature and sea-ice concentration changes for all models in AMIP6, there is still considerable intermodel spread in polar warming contributions (Supplementary Figure S9). Consistent with Crook and Forster (2011), intermodel spread in polar ocean heat uptake outweighs intermodel spread in most polar feedbacks for this modelled historical period, while intermodel spread in the albedo feedback plays a relatively larger role under CO₂ quadrupling.

**DISCUSSION**

Analysis of polar warming in CMIP6 reveals key contributors to polar amplification and their changes from CMIP5. While CMIP6 models overestimate historical Antarctic warming, they generally capture the observed pattern of strong Arctic amplification and weaker Antarctic warming. As in reanalysis data, Arctic warming in CMIP6 models is both surface- and winter-amplified, although CMIP6 shows stronger mid-tropospheric warming than the ERA-Interim reanalysis and previous climate models.

Our quantification of contributions to polar warming in CMIP6 is largely consistent with previous results for CMIP5 (e.g., Pithan and Mauritsen, 2014; Goosse et al., 2018). As in CMIP5, abrupt CO₂ quadrupling experiments in CMIP6 demonstrate that the lapse-rate and albedo feedbacks are the largest contributors to both Arctic amplification and weaker warming in the Antarctic than Arctic. The albedo feedback also contributes most to intermodel spread in polar warming, while the lapse-rate feedback and seasonal ocean heat storage contribute most to seasonal asymmetry in warming at both poles.

Novel results in comparison to existing literature include our assessment of the sensitivity of polar warming contributions to the choice of radiative kernel. While most feedbacks are relatively insensitive to kernel choice, the Arctic albedo warming contribution in CMIP6 varies by almost a factor of two for different kernels. This yields an Arctic albedo warming contribution in CMIP6 of equal or greater importance than the lapse-rate feedback in polar warming for half of the kernels considered, while the other half suggest that the lapse-rate feedback contributes more to polar amplification. However, the kernels most consistent with observations produce a stronger Arctic albedo feedback than previously found, on par with the lapse-rate feedback in its contribution to Arctic amplification. We also add a partitioning of AHT changes into moist and dry components, which demonstrates that increased moist AHT contributes to stronger Arctic amplification, and is the largest contributor to warming in the Antarctic. We find that increased polar warming in CMIP6 versus CMIP5 is explained by a stronger albedo feedback at both poles, combined with a less-negative cloud feedback in the Arctic and a larger increase in moist AHT to the Antarctic. Lastly, similar factors contribute to historical Arctic amplification in AMIP6 models compared to CMIP6 CO₂-quadrupling experiments, although the albedo feedback plays a larger role in weakening Antarctic warming in AMIP6 compared to CMIP6.

A limitation of using warming contribution methods to diagnose the mechanisms of polar amplification is that it implicitly includes interactions between feedbacks, making mechanistic interpretation difficult. For example, the strength of the lapse-rate feedback may be impacted by the amount of surface warming contributed by the albedo feedback (Graversen et al., 2014; Feldl et al., 2017) and mixed-phase cloud changes (Tan and Storelvmo, 2019), but the warming contribution method diagnoses the contributions of surface albedo, cloud, and lapse-rate changes separately. Others have argued that a strong winter lapse-rate feedback additionally requires seasonal ocean heat storage and sea-ice insulation loss in order to increase surface turbulent heat fluxes and upward longwave radiation, promoting warming in the lower-troposphere (Dai et al., 2019; Feldl et al., 2020; Chung et al., 2021). As demonstrated by these studies, experiments isolating specific mechanisms in climate models are needed to fully address the interconnected feedbacks promoting polar amplification.

While this study applies a TOA feedback framework, analysis from a surface perspective (e.g., Lu and Cai, 2009; Taylor et al., 2013; Pithan and Mauritsen, 2014; Boeke and Taylor, 2018) can provide additional insight on warming mechanisms in polar regions, where the surface and tropospheric responses to forcing are decoupled. Surface albedo and temperature feedbacks remain important contributors to Arctic amplification from a surface perspective in CMIP5 models, while the water vapor feedback and ocean heat uptake continue to oppose Arctic amplification (Pithan and Mauritsen, 2014). One key difference between surface and TOA frameworks in CMIP5 is that cloud feedbacks oppose Arctic amplification from a TOA perspective but slightly promote Arctic amplification from a surface perspective: changes in low-level clouds increase longwave radiation to the surface but have little impact on TOA longwave radiation, yielding a negative TOA cloud feedback due to shortwave radiative effects (Pithan and Mauritsen, 2014). Further analysis from a surface perspective may also be useful for understanding physical drivers of the surface-trapped warming that supports the positive Arctic lapse-rate feedback (Boeke et al., 2021).
than Arctic warming? While this may be explained by Clausius-Clapeyron nonlinearity and more-negative Antarctic cloud feedbacks, other possible mechanisms include any process (e.g., ocean heat uptake, water vapor feedback) that leads to a stronger equator-to-pole moist static energy gradient in the Southern Hemisphere than Northern Hemisphere under CO₂ quadrupling. Further investigation of these polar warming asymmetries may highlight key processes for constraining both Arctic and Antarctic amplification.

**DATA AVAILABILITY STATEMENT**

Publicly available datasets were analyzed in this study. This data can be found here: All CMIP and AMIP data analyzed for this study can be found in the Earth System Grid Federation (ESGF) repository at https://esgf-node.llnl.gov/projects/esgf-llnl/ . ERA-Interim data was provided by the ECMWF Data Archive at https://apps.ecmwf.int/datasets/data/interim-full-moda/ levtype=pl/. The HadCRUT5 Analysis is available from the University of East Anglia Climatic Research Unit at https://crudata.uea.ac.uk/cru/data/temperature/.

**AUTHOR CONTRIBUTIONS**

All authors contributed to the study design. MZ calculated the climate feedbacks and AMIP forcing, and LH computed the warming contributions and remaining analysis. LH wrote the original draft, and all authors contributed to the final manuscript.

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**SUPPLEMENTARY MATERIAL**

The Supplementary Material for this article can be found online at: https://www.frontiersin.org/articles/10.3389/feart.2021.710036/full#supplementary-material

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