Constrained Stabilization on the $n$–Sphere

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Abstract

We solve the stabilization problem on the $n$–sphere in the presence of conic constraints. We use the stereographic projection to map this problem to the classical navigation problem on $\mathbb{R}^n$ in the presence of spherical obstacles. As a consequence, any obstacle avoidance algorithm for navigation in the Euclidean space can be used to solve the given problem on the $n$–sphere.

We illustrate the effectiveness of the approach using the kinematics of the reduced attitude model on the $2$–sphere.

Key words: Constrained control; unit $n$–sphere; conic constraints, obstacle avoidance, Euclidean sphere world.

1 Introduction

Different mechanical systems of interest have state components that are restricted to evolve on the $n$–sphere. Example of such systems are the spherical pendulum (Shiriaev et al., 2004), the nonholonomic rolling sphere (Das and Mukherjee, 2004), the reduced-attitude or spin-axis stabilization of rigid bodies (Bullo and Murray, 1995; Tsiotras and Longuski, 1994) and the thrust-vector control for quad-rotor aircraft (Hua et al., 2009). Brockett (1973) developed a theory for the most elementary class of control problems defined on spheres where he discussed issues related to controllability, observability and optimal control without an explicit search for control laws. Bullo and Murray (1995) proposed a geometric approach to design controllers for control systems on the sphere relying on the notion of geodesics. The $n$–sphere is not diffeomorphic to a Euclidean space (it is a compact manifold without boundary) and, hence, there exist no smooth control law that globally stabilizes an equilibrium point on the $n$–sphere since the domain of asymptotic stability of any critical point of a continuous vector field needs to be diffeomorphic to a Euclidean space.

Recently, hybrid approaches have been proposed to guarantee global asymptotic and exponential stabilization on the $n$–sphere (Mayhew and Teel, 2013; Casau, 2015) and also on the group of rotations $\text{SO}(3)$ (Berkane and Tayebi, 2017; Berkane et al., 2017). However, in contrast to the unconstrained stabilization problem, there are only few research works that have considered the constrained stabilization problem on the $n$–sphere. This problem is relevant in different applications such as the pointing maneuver of a space telescope towards a given target (e.g., planets and galaxies) during which the telescope’s sensitive optical instruments must not be pointed towards bright objects such as stars. The control barrier functions (CBFs) approach on manifolds (Wu and Sreenath, 2015) can be used to solve the constrained stabilization problem on the $n$–sphere. However, this comes at the expense of solving a state-dependent online quadratic program and, besides, the domain of attraction is not characterized. For the particular case of the $2$–sphere, the constrained stabilization problem can be lifted to the constrained (full) attitude stabilization problem where different approaches exist. Spindler (2002) proposed a geometric control law that minimizes a given cost function to solve the problem of maneuvering a rigid spacecraft attitude from rest to rest while avoiding a single forbidden direction. In (Lee and Mesbahi, 2014), a logarithmic barrier potential function is used to synthesize a quaternion-based feedback controller that solves the attitude reorienta-
tion of a rigid body spacecraft in the presence of multiple attitude-constrained zones. Another potential-based approach for the constrained attitude control on $SO(3)$ has been proposed in (Kuhurni and Lee, 2017).

In this work we consider the constrained stabilization problem of dynamical systems evolving on a configuration space defined by the unit $n$–sphere. The considered forbidden zones are conic-type constraints in the sense that we force the state on the $n$–sphere (which can be seen as a unit axis) to keep minimum safety angles with respect to some given set of other unit axes. Our proposed solution consists first in showing that the considered constrained $n$–sphere manifold is diffeomorphic, via the stereographic projection (Helmske and B. Moore, 1996), to a Euclidean space punctured by spherical obstacles. Then, by considering a generic driftless system on the $n$–sphere, we prove that the pushforward vector field in the new stereographic coordinates is feedback linearizable. Therefore, we are able to map the given constrained stabilization problem on the $n$–sphere to the well-established and treated obstacle avoidance problem in $\mathbb{R}^n$ which allows us to benefit from the many studies in the latter field. For instance, one can use navigation functions (Koditschek and Rimon, 1990) to obtain an almost global result or even global results with hybrid control techniques (Berkane et al., 2019). We show that the qualitative properties (e.g., invariance, stability, region of attraction) of any static obstacle avoidance controller on $\mathbb{R}^n$ are preserved for the resulting safety controller on the $n$–sphere. **Notation:** We use $\mathbb{N}$, $\mathbb{R}$ and $\mathbb{R}_{\geq 0}$ to denote, respectively, the sets of positive integers, real and non-negative real numbers. $\mathbb{R}^n$ denotes the $n$–dimensional Euclidean space. $I_n$ denotes the $n \times n$ identity matrix and $e_k$ corresponds to the $k$–th column of $I_n$. The Euclidean norm of $x \in \mathbb{R}^n$ is defined as $\|x\| = \sqrt{x^\top x}$ where $(\cdot)^\top$ denotes the transpose of $(\cdot)$. The topological interior (resp. boundary) of a subset $S$ of a metric space is denoted by $\text{int}(S)$ (resp. $\partial S$). For a multi-variable function $f(x_1, \cdots, x_n)$, we denote by $\nabla f$ the gradient of $f$ with respect to the $i$–th argument $x_i$.

### 2 Problem Formulation

The unit $n$–sphere is an $n$–dimensional manifold that is embedded in the Euclidean space $\mathbb{R}^{n+1}$ and defined as $S^n := \{x \in \mathbb{R}^{n+1} : \|x\| = 1\}$. The tangent space to $S^n$ at a given point $x$ is defined by the $n$–dimensional hyperplane $T_x S^n = \{x \in \mathbb{R}^{n+1} : z^\top x = 0\}$, which represents all vectors in $\mathbb{R}^{n+1}$ that are perpendicular to $x$ in $S^n$. $S^n$ is a metric space if we pair it with the geodesic distance $d(x, y) := \arccos(x^\top y), \forall x, y \in S^n$. We consider the following driftless system on $S^n$:

$$\dot{x} = \Pi(x)u,$$  \hspace{1cm} (1)

where $u \in \mathbb{R}^m$ is the control input and $\Pi : \mathbb{R}^{n+1} \rightarrow \mathbb{R}^{(n+1) \times m}$ is a smooth matrix-valued function such that $\text{Im}(\Pi(x)) \subseteq T_x S^n$. The condition $\text{Im}(\Pi(x)) \subseteq T_x S^n$ implies that $\dot{x} \in T_x S^n$ or $x^\top \dot{x} = 0$ which guarantees forward invariance of $S^n$ under the dynamics (1) since $\|x\|^2$ remains constant regardless of the input $u$. Our goal is to propose a constrained stabilization strategy on $S^n$ in the presence of the following ($I + 1$) conic constraints:

$$O_i = \{x \in S^n : x^\top a_i > \cos(\theta_i)\}, i \in I = \{0, \cdots , I\} \hspace{1cm} (2)$$

where $a_i \in \mathbb{R}^n$ is the center of $O_i$ and $\theta_i \in (0, \pi/2)$ is the smallest angle (between $x$ and $a_i$) allowed in the free region. We define our constrained space on $S^n$ as $\mathcal{M} := S^n \setminus \bigcup_{i \in I}O_i = \{x \in S^n : d(x, a_i) \geq \theta_i, \forall i \in I\}$.

**Assumption 1** The following assumptions hold:

1. For all $x \in S^n \setminus \{e_{n+1}\}$, $\text{rank}(\Pi(x)) = n$.
2. For all $i, i' \in I$ with $i \neq i'$, $a_i a_{i'} < \cos(\theta_i + \theta_{i'})$.
3. $a_0 = e_{n+1}$.
4. $x(0) \in \mathcal{M}$ and $x_d \in \text{int}(\mathcal{M})$.

Item (1) of Assumption 1 is a controllability assumption that imposes the fact that we can steer any point on $S^n \setminus \{e_{n+1}\}$ in any direction by appropriately choosing the control input $u$. Item (2) imposes that the closures of the constraint zones $O_i$ are pairwise disjoint. In item (3) we assume, without loss of generality, that the obstacle $O_0$’s axis coincides with the coordinate axis $e_{n+1}$. Finally, item (4) imposes that the initial condition $x(0)$ and the desired reference point $x_d$ must lie in the free space $\mathcal{M}$ and the interior of $\mathcal{M}$, respectively.

**Problem 1** For system (1) and under Assumption 1, design a control law $u = k(x, x_d)$ that renders the constrained space $\mathcal{M}$ forward invariant and the target point $x = x_d$ on asymptotically (or exponentially) stable equilibrium with a region of attraction $R(x_d) \subseteq \mathcal{M}$.

### 3 Main Results

The stereographic projection is defined by the map $\psi : S^n \setminus \{e_{n+1}\} \rightarrow \mathbb{R}^n$ (Helmske and B. Moore, 1996)

$$\psi(x) = \begin{bmatrix} I_n & 0_{n \times 1} \end{bmatrix} x = \frac{J_n x}{1 - e_{n+1}^\top x} \hspace{1cm} (3)$$

Geometrically speaking, the stereographic projection of a point $x \in S^n \setminus \{e_{n+1}\}$ represents the unique point $\psi(x)$ describing the intersection of the line, that passes by $e_{n+1}$ and $x$, with the hyperplane $\{x \in \mathbb{R}^{n+1} : x_{n+1} = 0\}$. The following are some useful properties of this map.

**Lemma 1** The stereographic projection map $\psi$ satisfies:

- $\psi$ is a diffeomorphism with the inverse given explicitly by the map $\psi^{-1} : \mathbb{R}^n \rightarrow S^n \setminus \{e_{n+1}\}$ such that

$$\psi^{-1}(\xi) := \frac{2J_n \xi + (\|\xi\|^2 - 1)e_{n+1}}{1 + \|\xi\|^2} \hspace{1cm} (4)$$

The map $\psi^{-1}$ can be interpreted as the intersection of the line $\{x \in \mathbb{R}^{n+1} : x \|^2 = 1\}$ with $S^n$ and $\mathcal{M}$.
The Jacobian matrix of \( \psi(x) \) is given by
\[
\nabla \psi(x) = J_n \left( (1 - e\top_{n+1}x)I_{n+1} + xe\top_{n+1} \right) (1 - e\top_{n+1}x)^2 .
\]

For all \( x, x' \in S^n \setminus \{ e_{n+1} \} \)
\[
\| \psi(x) - \psi(x') \|_2^2 = \frac{2(1 - x\top_1x_2)}{(1 - e\top_{n+1}x_1)(1 - e\top_{n+1}x_2)} .
\]

If item (1) of Assumption 1 holds, then \( \Sigma(x) \) is full row rank (right invertible) and, hence, its Moore-Penrose

PROOF. The explicit expression \( \psi^{-1} \) is taken from (Helmke and B. Moore, 1996, Appendix C.4). Also, both \( \psi \) and \( \psi^{-1} \) are differentiable on their domains of definition and, hence, \( \psi \) is a diffeomorphism. The Jacobian of \( \psi \) is obtained by direct differentiation of (3). Making use of \( J_n^t J_n = I_{n+1} - e_{n+1}e_{n+1}\top \), we have

\[
\| \psi(x) \|_2^2 = \frac{1 - (e_{n+1}x)^2}{(1 - e_{n+1}x)^2} = \frac{1 + e_{n+1}x}{1 - e_{n+1}x} .
\]

Therefore, for all \( x, x' \in S^n \setminus \{ e_{n+1} \} \)
\[
\| \psi(x) - \psi(x') \|_2^2 = \| \psi(x) \|_2^2 - 2 \psi(x)\top \psi(x') + \| \psi(x') \|_2^2
\]
\[
= \frac{1 + e_{n+1}x_1 + e_{n+1}x_2}{1 - e_{n+1}x_1} - 2 \frac{2x_1^\top J_n x_2}{1 - e_{n+1}x_1} + \frac{1 + e_{n+1}x_2}{1 - e_{n+1}x_2} - 2 \frac{2x_1^\top J_n x_2}{1 - e_{n+1}x_2}
\]
\[
= \frac{2(1 - x_1^\top x_2)}{(1 - e_{n+1}x_1)(1 - e_{n+1}x_2)} .
\]

Since \( \psi \) is a diffeomorphism, we have \( \forall x \in S^n \setminus \{ e_{n+1} \} \), rank(\( \nabla \psi(x) \)) = \( n \). However, \( \nabla \psi(x)(x - e_{n+1}) = 0 \) and thus ker(\( \nabla \psi(x) \)) = \{ \( \alpha(x - e_{n+1}) : \alpha \in \mathbb{R} \) \}. On the other hand, since \( x \in S^n \setminus \{ e_{n+1} \} \), we have \( x\top (x - e_{n+1}) = 1 - x\top e_{n+1} \neq 0 \) which implies that \( \alpha(x - e_{n+1}) \notin T_{S^n} \) or ker(\( \nabla \psi(x) \)) \( \cap \) Im(\( \Pi(x) \)) = \( \emptyset \). Finally, by applying (Bernstein, 2009, Fact 2.10.14, item ii)), we have rank(\( \Sigma(x) \)) = rank(\( \Pi(x) \)) - dim(ker(\( \nabla \psi(x) \)) \( \cap \) Im(\( \Pi(x) \))) = rank(\( \Pi(x) \)) = \( n \).

We show in the following lemma that \( \psi \) maps the constrained space \( M \) to a Euclidean sphere world on \( \mathbb{R}^n \) as defined in (Koditschek and Rimon, 1990); see Fig. 1.

**Lemma 2** Let Assumption 1 hold. Then, the image of the constrained zones \( \{ O_i \}_{i \in I} \) via the map \( \psi \) is given by the following pairwise disjoint spherical subsets of \( \mathbb{R}^n \)
\[
\tilde{O}_0 := \psi(O_0) = \{ \xi \in \mathbb{R}^n : \| \xi \| > \cot(\theta_0/2) \},
\]
\[
\tilde{O}_i := \psi(O_i) = \{ \xi \in \mathbb{R}^n : \| \xi - c_i \| < r_i \}, \ i \in I \setminus \{ 0 \},
\]

\[
\tilde{O}_i := \psi(O_i) = \{ \xi \in \mathbb{R}^n : \| \xi - c_i \| = r_i \}, \ i \in I \setminus \{ 0 \}.
\]

The resulting Euclidean sphere world consists of one large \( (n - 1) \)-sphere \( \mathbb{R}^{n-1} \setminus \psi(O_0) \) that bounds the workspace and other smaller disjoint \( (n - 1) \)-spheres \( \psi(O_i) \) that define obstacles in \( \mathbb{R}^n \) that are strictly contained in the interior of the workspace. The obtained constrained space on \( \mathbb{R}^n \) is \( \tilde{M} := \psi(M) = \mathbb{R}^n \setminus \bigcup_{i \in I} \tilde{O}_i \). Now, let us consider the change of variable \( \xi := \psi(x) \). Then, in view of (1), the dynamics of \( \xi \) are given by
\[
\dot{\xi} = \nabla \psi(x)x = \nabla \psi(x)\Pi(x)u = \Sigma(x)u .
\]

Interestingly, according to Lemma 1, \( \Sigma(x) \) is full row rank (right invertible) and, hence, its Moore-Penrose
pseudo inverse can be explicitly calculated as follows:
\[ \Sigma(x)^+ = \Sigma(x)^T (\Sigma(x)\Sigma(x)^T)^{-1}. \tag{10} \]
Therefore, by considering a control law of the form
\[ u = \Sigma(x)^+ v, \tag{11} \]
where \( v \in \mathbb{R}^n \) is a virtual control input, one obtains
\[ \xi = v. \tag{12} \]

Next, we show in Theorem 1 that solving Problem 1 on the constrained space \( \mathcal{M} \) boils down to solving the following problem on the Euclidean sphere \( \mathbb{M} \).

**Problem 2** For system (12) and for a given \( \xi_d \in \text{int}(\mathcal{M}) \), design a control law \( v = \hat{\kappa}(\xi, \xi_d) \) that renders the constrained space \( \mathcal{M} \) forward invariant and the target point \( \xi = \xi_d \) an asymptotically (exponentially) stable equilibrium with a region of attraction \( \mathcal{R}(\xi_d) \subseteq \mathcal{M} \).

**Theorem 1** If \( v = \hat{\kappa}(\xi, \xi_d) \) is a control law that solves Problem 2 with region of attraction \( \mathcal{R}(\xi_d) \) then \( u = \Sigma(x)^+ \hat{\kappa}(\psi(x), \psi(x_d)) \) is a control law that solves Problem 1 with a region of attraction \( \mathcal{R}(x_d) = \psi^{-1}(\mathcal{R}(\psi(x_d))) \).

Moreover, if \( \hat{\kappa}(\xi, \xi_d) \) is a priori bounded on \( \mathcal{M} \) then the control law \( u \) is also a priori bounded on \( \mathcal{M} \).

**PROOF.** To prove this result we use the comparison theorem (Michel et al., 2001, Theorem 3.4.1) to deduce the qualitative properties of all solutions of
\[ \dot{x} = \Pi(x)\Sigma(x)^+ \hat{\kappa}(\psi(x), \psi(x_d)), x(0) \in \mathcal{M}, \tag{13} \]
from the qualitative properties of solutions of
\[ \dot{\xi} = \hat{\kappa}(\xi, \xi_d), \xi(0) = \psi(x(0)), \xi_d = \psi(x_d). \tag{14} \]
We denote hereafter by \( \mathcal{S}_x \) (resp. \( \mathcal{S}_\xi \)) the set of all solutions to (13) (resp. (14)). It is clear that \( \psi(\mathcal{S}_x) \subseteq \mathcal{S}_\xi \) since, for all \( p_x(t, x) \in \mathcal{S}_x \), we have \( \psi(p_x) = \nabla \psi(p_x)p_x = \Sigma(p_x)\Sigma(p_x)^+ \hat{\kappa}(\psi(p_x), \psi(x_d)) = \hat{\kappa}(\psi(p_x), \xi_d) \) and hence \( \psi(p_x(t, x)) \in \mathcal{S}_\xi \). Moreover, in view of (6) we have
\[
\|\psi(x) - \psi(x_d)\|^2 = \frac{2(1 - x^T x_d)}{(1 - e_{n+1}^T x)(1 - e_{n+1}^T x_d)} \geq \frac{2(1 - \cos(d(x, x_d)))}{(1 - \cos(d(x, e_{n+1})))(1 - \cos(d(x, e_{n+1})))} \geq \frac{2\sin^2(d(x, x_d))}{2}\frac{\sin^2(d(x, e_{n+1}))/2}{\sin^2(d(x, e_{n+1}))/2}. \tag{15}\]
Now, using the fact that \( x, x_d \in \mathcal{M} \subseteq \mathbb{R}^n \setminus \mathcal{O}_0 \) (i.e., \( d(x, e_{n+1}) \geq \theta_0 \) and \( d(x_d, e_{n+1}) \geq \theta_0 \)) and the useful identity \((2/\pi)z \leq \sin(z) \leq z, z \in [0, \pi/2]\), one deduces
\[
4\pi^2 d^2(x, x_d) \leq \|\psi(x) - \psi(x_d)\|^2 \leq \sin^{-4}(\theta_0) d^2(x, x_d). \tag{16}\]
Since the distance \( \|\psi(x) - \psi(x_d)\| \) on the closed set \( \mathcal{M} \) is upper and lower bounded by class \( \mathcal{K}_\infty \) functions (more precisely positive-power functions) of the distance \( d(x, x_d) \) on the closed set \( \mathcal{M} \), we can apply (Michel et al., 2001, Theorem 3.4.1) to conclude that forward invariance of \( \mathcal{M} \) and \( \xi = \xi_d \) (with respect to (14)) implies, respectively, forward invariance of \( \mathcal{M} \) and \( x = x_d \) (with respect to (13)). Moreover, asymptotic (exponential) stability of the equilibrium \( \xi = \xi_d \) implies asymptotic (exponential) stability of the equilibrium \( x = x_d \).

In the case of a single constraint, one has \( \mathcal{M} = \mathbb{R}^n \setminus \mathcal{O}_0 \) which, in view of Lemma 2, represents the ball bounded by the sphere of radius \( \cot(\theta_0/2) \) that is centered at 0. It is not difficult to show that setting \( \gamma = -\gamma(\xi - \xi_d) \) in (12) results in GES of \( \xi = \xi_d \) and forward invariance of \( \mathcal{M} \). The following corollary follows from Theorem 1.

**Corollary 2 (single constraint)** Consider the kinematics (1) under Assumption 1 and let \( I = 0 \). Then, the control law \( u = -\gamma \Sigma(x)^+ (\psi(x) - \psi(x_d)), \gamma > 0 \), guarantees global exponential stability of the equilibrium \( x = x_d \) and forward invariance of the free space \( \mathcal{M} \).

The global result of Corollary 2 is related to the well-known Alexandroff one-point compactification in general topology (Alexandroff, 1924). In fact, removing a single constraint zone \( \mathcal{O}_0 \) from the unit \( n \)-sphere results in a manifold that is diffeomorphic to a Euclidean space and, therefore, global asymptotic stability is possible via a continuous time-invariant feedback.

If we have two or more constraints, the constrained manifold \( \mathcal{M} \) is not diffeomorphic to a Euclidean space and, hence, there is a topological obstruction to solve Problem 2 globally with a continuous feedback. Different controllers from the vast literature on obstacle avoidance.
can be employed here. For instance, continuous feedback, e.g., (Koditschek and Rimon, 1990; Loizou, 2017), can be used to ensure almost global asymptotic stabilization while hybrid feedback, e.g., (Berkane et al., 2019), can be used to ensure stronger global asymptotic stabilization. In this work, we consider the navigation functions-based approach of (Koditschek and Rimon, 1990). In particular, let the navigation function

$$\phi(\xi, \xi_d) = \frac{||\xi - \xi_d||^2}{(||\xi - \xi_d||^{2k} + \beta(\xi))^{\frac{k}{2}}}$$ \quad k > 0, \quad (18)$$

where $$\beta(\xi) = \Pi_{i \in I} \beta_i(\xi)$$ and

$$\beta_0(\xi) = \cot^2(\theta_0/2) - ||\xi||^2, \quad \beta_i(\xi) = ||\xi - c_i||^2 - r_i^2, \quad i \in I \setminus \{0\}.$$ \quad (19, 20)

Note that the parameter $$k$$ needs to be tuned above a certain threshold in order to eliminate local minima and for $$\phi$$ to be a valid navigation function. We then consider navigation along the negative gradient of $$\phi$$ and define

$$\ddot{\kappa}(\xi, \xi_d) := -\nabla_\xi \phi(\xi, \xi_d), \quad \gamma > 0.$$ \quad (21)

We state the following corollary that follows from the result of Theorem 1 and (Koditschek and Rimon, 1990).

**Corollary 3 (Two or more constraints)** Consider the kinematics (1) under Assumption 1. Then, there exists $$K$$ such that if $$k > K$$ the control law $$u = \Sigma(\xi)^0 \ddot{\kappa}(\psi(\xi), \psi(\xi_d))$$, with $$\ddot{\kappa}$$ defined in (21), guarantees almost global asymptotic stability of the equilibrium $$x = x_d$$ and forward invariance of $$M$$.

**Remark 1** The feedback linearization approach in (9)–(12) can be extended to high-order dynamics of the form:

$$\begin{align*}
\dot{x}_1 &= \Pi(x_1)x_2, \\
\dot{x}_q &= x_{q+1}, 2 \leq q \leq (l - 1), \quad \dot{x}_l = u, \quad (22)
\end{align*}$$

where $$x = (x_1, x_2, \ldots, x_l) \in S^n \times \mathbb{R}^{(l-1)m}$$ and $$u \in \mathbb{R}^m$$. Let the change of variables $$\xi = (T_1(x), \ldots, T_l(x))$$ such that $$T_q(x) := \psi(x_1), T_2(x) := \Sigma(x_1)x_2$$ and

$$T_{q+1}(x) := \nabla_1 T_q(x)\Pi(x_1)x_2 + \sum_{p=2}^q \nabla_p T_q(x)x_{p+1}$$

for $$q = 2, \ldots, (l - 1)$$. By construction $$T_q(x)$$ depends only on $$x_1, \ldots, x_q$$ and, hence, $$\nabla_1 T_q + T_{q+1}(x) = \nabla_1 T_q(x)$$. It follows that $$\nabla_1 T_q(x) = \nabla_2 T_2(x) = \Sigma(x_1)$$ and the dynamics of the new variables are

$$\begin{align*}
\dot{\xi}_q &= \xi_{q+1}, 1 \leq q \leq l - 1, \\
\dot{\xi}_l &= \nabla_1 T_l(x)\Pi(x_1)x_2 + \sum_{p=2}^{l-1} \nabla_p T_l(x)x_{p+1} + \Sigma(x)u.
\end{align*}$$

Therefore, the control law

$$u = \Sigma(x_1)^0 (v - \nabla_1 T_l(x)\Pi(x_1)x_2 - \sum_{p=2}^{l-1} \nabla_p T_l(x)x_{p+1}),$$

where $$v \in \mathbb{R}^n$$, results in the linear dynamics

$$\dot{\xi}_q = \xi_{q+1}, 1 \leq q \leq l - 1, \quad \dot{\xi}_l = v.$$ \quad (23)

**4 Example**

We consider the kinematics of the spherical pendulum

$$\dot{x} = x \times u =: \Pi(x)u$$ \quad (24)$$

where $$\times$$ denotes the cross product and $$u$$ is the angular velocity of the pendulum. Using the cross product identities $$\Pi(x)^T = -\Pi(x)$$ and $$\Pi(x)^2 = -I_3 + xx^T$$, it is easy to show that $$\Sigma(x)^T = (1 - e_3 x)^{-2}I_2$$. It follows from (5) and (10) that $$\Sigma(x)^T = -\Pi(x)((1 - x_3)I_3 + e_3 x^T)I_2^T$$. For simulation, we pick $$x(0) = (-1, 0, 1)/\sqrt{2}, x_d = (1, 2, -2)/3$$ and $$\gamma = k = 5$$ for the control parameters. We consider 5 constraints zones such that $$a_0 = e_3, a_1 = e_1, a_2 = -e_1, a_3 = e_2$$ and $$a_4 = -e_2$$. The angles are given by $$\theta_i = \pi/(7 + i)$$ for all $$i = 0, 1, \ldots, 4$$. It is easy to check that Assumption 1 holds. Simulation results are plotted in Figure 2 which show a successful constrained stabilization on the unit 2–sphere in the presence of different constraint zones. The complete simulation video can be found at https://youtu.be/ye8delheiok.

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