Identification of Fracturing Behavior In Thermally Cracked Granite Using The Frequency Characteristics of Acoustic Emission

Kazumasa Sueyoshi (ksueyoshi@hiroshima-u.ac.jp)  
Hiroshima University  https://orcid.org/0000-0001-9976-4876

Manami Kitamura  
National Institute of Advanced Industrial Science and Technology

Xinglin Lei  
National Institute of Advanced Industrial Science and Technology

Ikuo Katayama  
Hiroshima University

Research Article

Keywords: Acoustic emission, peak frequency, thermal crack, Inada granite, triaxial compression test

Posted Date: December 8th, 2021

DOI: https://doi.org/10.21203/rs.3.rs-1126551/v1

License: This work is licensed under a Creative Commons Attribution 4.0 International License. 
Read Full License
Abstract

The frequency characteristics of acoustic emission (AE) during triaxial compression of thermally cracked and unheated (“fresh”) granite samples were investigated with the aim of understanding the influence of pre-existing cracks on precursor information regarding macroscopic failure. The peak frequency during the damage process was the same for thermally cracked and fresh granites. Analysis of AE signals showed that signals with low peak frequency appeared before failure of the sample, implying the initiation of microfractures with progressive growth of cracks. The peak amplitude of the frequency spectrum recorded in the thermally cracked samples was much lower than that in the fresh samples. This result suggests two reasons for the difference in peak amplitude: reduction in shear modulus and the attenuation filtering phenomenon caused by thermal cracks. In particular, the maximum value of peak amplitude in the low-frequency band for the thermally cracked samples was smaller than that for fresh samples. This characteristic can be related to the stress drop and crack size. Assuming that pre-existing thermal cracks grow during the pre-failure stage, the events with low peak frequency and low peak amplitude in the heat-treated samples are interpreted as exhibiting a low stress drop because of the small rupturing area for individual events. Therefore, although AE signals with low frequency can be considered as precursors to rock failure, cracking behavior suggested by events with low frequency depends on the initial damage condition of the rock sample.

Introduction

The physical and mechanical properties of rocks, which are influenced by cracks of varying size, are crucial for understanding the mechanism of hydraulic fracturing and induced seismicity related to many applied engineering fields, such as enhanced geothermal systems (EGSs), shale gas hydraulic fracturing, and disposal of wastewater from resource extraction (Fournier 1996; Rothert and Shapiro 2007). To analyze the influence of cracks on rock physical and mechanical properties, it is necessary to perform laboratory experiments on rocks with pre-existing cracks. In addition, it is also important to investigate the fracturing process starting from the occurrence of new cracks or propagation of pre-existing cracks through to fault formation. The effect of pre-existing cracks on rock physical and mechanical properties has been examined through laboratory experiments on thermally cracked samples. For example, Sun et al. (2015) performed uniaxial compression tests with granite specimens heated to 800 °C and showed that peak stress decreased as heating temperature increased. Wang et al. (2013) conducted hydrostatic tests at an effective pressure of up to 90 MPa and then carried out triaxial deformation tests at an effective pressure of 30 MPa with granite samples heat-treated at 700 °C. Those authors reported that, although dilatancy in the heat-treated samples started at lower differential stress compared with the non-heated samples, the thermal cracks had no influence on rock strength. From the observation that elastic wave anisotropy appeared at an early stage of triaxial deformation of the thermally cracked samples, the apparent insensitivity of rock strength to thermal cracks was interpreted as being due to stress shielding induced by aligned crack distributions. Moreover, microstructural analysis revealed the existence of long cracks in both non-heated and heat-treated samples, which was interpreted in terms of the more effective
control of rock strength by such pre-existing cracks compared with thermal cracks, which are shorter. Therefore, it would appear that measuring the effect of pre-existing cracks is essential for estimating the strength of rocks.

In elastic rocks, the creation of new cracks and growth of pre-existing cracks cause dynamic energy release, which propagates in the form of an elastic wave and is termed “acoustic emission” (AE). AE monitoring is an effective approach for understanding the fracturing process in rocks (e.g. Lei and Ma 2014). At the laboratory scale, deformation processes have been inferred from the AE event rate and the spatial distribution of AE hypocenters (e.g., Scholz 1968; Lockner et al. 1991; Lei et al. 2000; Zang et al. 2000; Thompson et al. 2005). AE parameters such as frequency and amplitude are related to the length of developed cracks and have been used to interpret precursor information regarding macroscopic failure. Wang et al. (2019) studied the spectral characteristics and distribution of the energy ratio of the AE signals band acquired by fast Fourier transform (FFT) and wavelet packet transform (WPT), and showed that the low-frequency band of the AE signal represented the formation of macrocracks. Chen et al. (2020) investigated the relationship between AE frequency spectrum characteristics and cracking behaviors using high-speed camera images. Their results revealed that the signal with high peak amplitude in low-frequency bands can be used as precursor information regarding macroscopic failure. However, the relationship between AE frequency spectrum characteristics and pre-failure damage has not yet been explained in detail.

The present study aimed to confirm the influence of pre-existing cracks on AE frequency spectrum characteristics and cracking activity. We measured AE signals during triaxial compression of both thermally cracked and unheated granite samples under conditions of constant confining pressure and temperature. Then, AE frequency spectra were extracted from each signal. The effect of pre-existing cracks on precursor information regarding macroscopic failure was ascertained by comparing differences in precursory features between thermally cracked and non-heated samples. By elucidating the precursors of failure, it is possible to predict the failure of rocks on an experimental basis and in turn to perform risk assessment of induced seismicity in EGSSs, as well as in other domains in which artificial cracks must be created, such as hydraulic fracturing of reservoirs.

Method

Samples

For this study, we used samples of Inada granite from Ibaraki, Japan, which is typically composed of 36 vol.% quartz, 32 vol.% plagioclase, 28 vol.% alkali-feldspar, and 4 vol.% biotite, with an average grain size of 2–4 mm (Lin 2002). Inada granite is structurally anisotropic as a result of its partially oriented arrangement of mineral particles and pre-existing microcracks, and the anisotropy of its compressive strength is approximately 10% (Lin et al. 2003). All samples used in our experiments were shaped into cylinders of 50 mm diameter and 125 mm length and were cored perpendicular to the plane that had the highest crack density so as to reduce the effect of anisotropy on the deformation process. P-wave
velocity was measured in the directions parallel ($V_{\text{axial}}$) and perpendicular ($V_{\text{lateral}}$) to the axis of the sample. $V_{\text{lateral}}(\text{MAX})$ and $V_{\text{lateral}}(\text{MIN})$ were determined as the maximum and minimum values of P-wave velocity measured in three directions every 60° at the mid-point of the sample length, respectively. The velocity difference $V_{\text{lateral}}(\text{MAX}) - V_{\text{axial}}$ was 0.263–0.693 km/s at room conditions.

Inada granite samples (IG2 to IG5) were heated in an oven at room pressure, up to 550°C, at an average rate of 52°C/min, after which the temperature was kept constant at 550°C for 45 min, and finally cooled slowly overnight in the turned-off oven. Cracking occurred in the samples as a result of thermal stressing due to the differing thermal expansion of neighboring grains. In this paper, we refer to heat-treated samples as “IGh” and fresh samples as “IGf”. Microscopic observations show that the heating treatment led to the formation of wide cracks ranged to 100 µm (Fig. 1). The initial porosities $\phi$ of the fresh and thermally cracked samples were 0.41–0.92% and 1.59–1.81%, respectively (Table 1). Porosity was determined by

$$\phi = \left( \frac{W_{\text{wet}} - W_{\text{dry}}}{\rho_{\text{water}}V_{\text{bulk}}} \right)$$

where $W_{\text{wet}}$ and $W_{\text{dry}}$ are the water-saturated and dry weights of the sample, respectively, $\rho_{\text{water}}$ is the density of water, and $V_{\text{bulk}}$ is the bulk volume of the sample, which can be determined by the cross-sectional area and length of the sample. The absolute error on porosity measurements for our samples is 0.00575–0.00587%, based on the uncertainty of the size and weight values.

**Experimental Setup**

Triaxial compression tests were conducted using the loading system at the Geological Survey of Japan, National Institute of Advanced Industrial Science and Technology, Ibaraki, Japan (Fig. 2a) (Lei 2006). The loading system was renewed recently, allowing the axial piston to be advanced at a constant stress/strain rate using a servo-controlled system. During the experiments, the confining and pore pressures were controlled individually by syringe pumps. Axial ($\varepsilon_a$) and radial ($\varepsilon_r$) strains were measured with six cross-type strain gauges attached directly to the samples (Fig. 2b). Twenty-eight piezoelectric transducers (PZTs) with resonance frequencies of 0.25–2 MHz were mounted on the sample surface for monitoring of AE and measurement of velocity (Table 2). After all of the strain gauges and PZTs had been glued onto the samples, the specimens were jacketed with silicone sealant (about 5 mm in thickness) to prevent interaction with the confining oil. All experiments were performed at a constant displacement rate of 0.01 mm/min (corresponding to a strain rate of $\sim1.33 \times 10^{-6}$ s$^{-1}$) and a constant confining pressure of 24.5 MPa. For wet experiments, distilled water was used as pore fluid and injected into the upstream of the sample at a constant pore pressure of 4 MPa. The downstream of the sample was released to atmospheric pressure.

A total of five samples were tested. In this study, we focused on results for two samples: IGh5 and IGf8. Samples IGh5 and IGf8 were selected to investigate the frequency characteristics of AE signals because a relatively wide range of dominant frequency bands can be measured using PZTs with resonant frequencies of 2 and 0.25 MHz mounted on the two samples. Samples IGh2, IGh3, and IGf6 were chosen for checking differences in frequency analysis results depending on the resonant frequency of PZTs.
Data Processing

In the experiments, volumetric strain \( (\varepsilon_v) \) was calculated from \( \varepsilon_a \) and \( \varepsilon_r \) using \( \varepsilon_v = \varepsilon_a + 2\varepsilon_r \). All strain values were averaged from six strain gauges. During the initial stages of deformation, volumetric strain increases linearly with applied differential stress. The result for triaxial compression testing of thermally cracked Inada granites (IGh2) show that the samples underwent the pre-existing crack closure stage at \( \sim 15\% \) of the peak stress. Therefore, the volumetric strain curve between 15\% and 25\% of the peak stress is presumed to be a linear region (i.e., one characterized by a linear relationship). In this study, the stress at which volumetric strain deviates from the linear region is defined as the onset of dilatancy. At high differential stress, unstable crack growth, which is termed “crack damage stress”, starts at a stress level between 70\% and 85\% of the peak stress, and this point corresponds to the maximum volumetric strain in the strain curve (Bieniawski 1967; Martin and Chandler 1994). In this study, the differential stress at the maximum volumetric strain is defined as crack damage stress.

AE events were recorded automatically by a fast waveform recording system, as described by Lei (2006). The signals detected from every PZT were pre-amplified by 40 dB and then fed into the waveform system. We digitized the AE signals with a sampling rate of 25 MHz and a sample length of 4096 words (~160 µs). AE waveform data were acquired by the following triggering logic. First, four transducers were selected as triggering sensors (p10, 13, 21, and 24 in sample IGh2-3 and p9, 10, 23, and 24 in samples IGh5, IGf6, and IGf8). These transducers were distributed homogeneously on the specimen surface to avoid deviation of triggering. An OR logic (triggering when at least one of selected inputs overshoots the threshold level) was used within the four transducers, which enabled us to record AE waveforms when any one of the transducers detected a signal exceeding the pre-defined threshold level (50 mV in this study). Triggered AE events were recorded using single-event mode, in which the waveform data were immediately transferred to the hard disk of the host PC and for which approximately two events per second can be recorded during the entire loading process. In the experiment involving sample IGf8, after the AE event rate increased, the system was switched to multi-event mode, in which approximately 5000 events per second can be recorded in the on-board memory. Thus, the number of AE events recorded for IGf8 was not comparable with other experiments. To investigate the relationship between cracking behavior and AE peak frequency, the frequency spectrum was obtained using FFT, as follows:

\[
A(f) = \int_{-\infty}^{\infty} A(t) e^{-i2\pi ft} dt
\]

where \( A(t) \) is the original AE waveform, \( A(f) \) is the frequency spectrum, \( f \) is the frequency component of the signal, and \( t \) is the acquisition time of \( A(t) \). Using Equation (1), we utilized the original AE waveform (Fig. 3a) to obtain the frequency spectrum (Fig. 3b). In the frequency spectrum, the peak amplitude is the largest amplitude in the spectrum, and the peak frequency is the point at the peak amplitude (Fig. 3b). The transducer that recorded the most events was selected, and then the frequency spectrum characteristics were extracted from the signals recorded by the selected transducers. In the experiment involving sample IGh3, sensors p10 and p13 were chosen as the typical transducers with resonant
frequencies of 1 and 0.5 MHz, respectively, and in the experiment involving IGh5, sensors p23 and p9 were selected as the representative instruments, with resonant frequencies of 2 and 0.25 MHz, respectively. In the case of fresh samples IGf6 and IGf8, p10 and p9 were used for frequency analysis as the typical sensors, with resonant frequencies of 1 or 2 MHz and 0.5 or 0.25 MHz, respectively. In addition, we selected the transducer close to a region where AE events are densely distributed (p3) and obtained the reanalyzed peak frequency from the signal detected by the selected transducer (detailed in section 3.3).

In addition to the peak frequency, we also estimated the corner frequency of the AE events. Crack models suggest that the P or S displacement spectrum is flat at low frequencies and decreases proportionally to $f^{-2}$ at high frequencies (Aki 1967; Sato and Hirasawa 1973). The high- and low-frequency trends intersect at the so-called corner frequency, which is the key parameter used for estimating fault size and stress drop in seismology. The AE spectrum calibrated by the average values of the difference between the power spectral densities of PZT and laser Doppler vibrometer (LDV) records is not flat and is roughly proportional to $f^2$ in the low-frequency domain (Lei et al. 2003). However, as the frequency response of PZTs was not calibrated using LDV in our experiments, it is difficult to consider the intersection between the fitting lines of $f^2$ and $f^{-2}$ as the corner frequency. In this study, the frequency of the point marking the largest amplitude in the spectrum fitted by an autoregressive (AR) model (Akaike, 1998) was obtained as the corner frequency (Fig. 3c) and averaged over all sensors. The duration of the signal used for the spectral analysis was 20 µs from 5 µs before the first motion. The spectral analysis was performed on waveforms recorded by all PZTs, and the median value was extracted as the typical corner frequency at each event.

Experimental Results

Mechanical data and AE event rate

The maximum differential stress of thermally cracked granite under wet conditions ranged from 357 to 413 MPa, whereas that of fresh granite under wet conditions was 431 MPa (Table 2). After reaching peak stress, brittle fracture occurred rapidly in all of the experiments. Results for the thermally cracked samples indicate that dilatancy began at 27.1–32.7% of the peak stress and that crack damage stress started at 50.4–62.6% (Fig. 4a–c). For the fresh samples, dilatancy began at 29.8–36.7% and crack damage stress started at 58.2–68.5% of the peak stress (Fig. 4d, e). Systematic trends between the thermally cracked and fresh samples in strain behavior were not observed. Figure 4f–j shows differential stress and AE events rate obtained from the waveform recording system as functions of elapsed time. AE data obtained from the wet samples show that cracking activity was initiated at a stress level of ~85% of the peak stress in the sample using PZTs with a resonant frequency of 1 MHz (IGh2) (Fig. 4f). In contrast, the AE event rate began to increase at around the crack damage stress in the sample using PZTs with resonant frequencies of 1 and 0.5 MHz (Fig. 4g, IGh3). A similar behavior was observed for the sample using PZTs with resonant frequencies of 2 and 0.25 MHz (IGh5). From the above results, even in the experiments
under wet conditions, it was possible to record a small amount of AE at the point of crack damage stress by adding PZTs with low-frequency bands (0.5 or 0.25 MHz). AE events in the fresh sample were successfully obtained by using PZTs with resonant frequencies of 1 and 0.5 MHz (Fig. 4i). After the differential stress reached ~50% of the expected peak stress in IGf8, the system was switched to multi-mode, meaning that the AE events rate increased rapidly from ~4 minutes before failure (Fig. 4j). This paper does not consider the number of AE counts in experiment IGf8.

**Peak Frequency**

The peak frequency of each AE event as a function of time for thermally cracked granites is shown in Figure 5. In all experiments, the peak frequency was distributed mainly in the range of 0–0.6 MHz. High-peak-frequency AE signals ranging from 0.6 to 2 MHz were detected only by PZTs with a high resonant frequency of 2 MHz. Focusing on the range of 0–0.6 MHz, the AE signals can be divided into two peak frequency bands. The peak frequency bands in IGh3 were distributed in similar bands: a low-frequency band of 0.1–0.25 and a high-frequency band of 0.25–0.45 MHz for the PZTs with resonant frequencies of 0.5 and 1 MHz (Fig. 5a, b). Results for PZTs with resonant frequencies of 0.5 and 1 MHz reveal that events with intermediate peak frequency of 0.2–0.3 MHz increased with increasing differential stress. In contrast, AE events recorded by the PZT with a resonant frequency of 0.25 MHz in IGh5 were divided into different bands: a low-frequency band of 0.05–0.15 and a high-frequency band of 0.2–0.3 MHz in the PZT with a resonant frequency of 0.25 MHz (Fig. 5c). As shown in Figure 5d, AE signals with lower frequencies were measured by the PZT with a resonant frequency of 2 MHz as differential stress was increased. The peak frequency of AE events in the fresh granites shows a similar distribution to that of events in the thermally cracked granites (Fig. 6). Signals with the resonant frequency of the PZT and its integer multiple frequencies are likely to be reflected in the frequency spectrum, and the frequency response of PZT sensors is not flat. Thus, the frequency bands of the peak frequency that resulted from the response of PZTs to harmonic waves are included in the AE signals. A PZT sensor with a resonant frequency of 2 MHz has a relatively uniform frequency response in the range of 0 to 0.6 MHz, as revealed by experiments (Lei et al., 2003). Thus, it is probable that the appearance of the events with lower frequencies recorded by the PZT with 2 MHz resonant frequency indicates the frequency characteristics that are independent of the resonant frequency of the PZT. The frequency characteristics recorded by PZTs with a 2 MHz resonant frequency are discussed in detail in the section “Discussion”.

**Dominant Frequency**

To clarify the relationship between the dominant frequency and cracking behavior, the peak frequency for PZTs with a resonant frequency of 2 MHz was compared with the corner frequency. After identifying a region where AE events are densely distributed (the center of the sample in this study) and a transducer with a small incidence angle and short distance from the region, we analyzed the waveform that was radiated from the region and recorded by the transducer. The peak frequency extracted by this process
was designated as the “reanalyzed peak frequency”. Figure 7 shows the reanalyzed peak frequency and corner frequency as functions of time for both thermally cracked and fresh samples. The reanalyzed peak frequency was 0.4–0.5 MHz in the middle stage of loading and ranged from 0.1 to 0.5 MHz at the time close to failure. The tendency for the number of AE events dominated by low-frequency components to increase with loading can be observed from the results for corner frequency. The magnitude of the dominant frequency of AE is related to the source dimension (crack size). According to Sato and Hirasawa (1973), by assuming that the AE signal is radiated from the circular crack that grows at a constant velocity, the source dimension $L$ can be expressed by the corner frequency of a P-wave, $f_0(P)$, as

$$L = \frac{C_p V_p}{2\pi f_0(P)},$$  \hspace{1cm} (2)

where $C_p$ is an implicit function of the rupture velocity as presented in Sato and Hirasawa (1973) and $V_p$ is the P-wave velocity. The relationship expressed in Equation (2) indicates that the source dimension is inversely proportional to the corner frequency. Lei et al. (2003) reported that a corner frequency of 0.1 MHz corresponds roughly to a source dimension of 10–15 mm. As the corner frequency and reanalyzed peak frequency obtained in the present study have the same tendency during loading, it is considered that the increase in the number of AE events with low peak frequency with increasing differential stress represents the formation of macrocracks.

**Discussion**

**Interpretation of AE frequency spectrum characteristics**

For the thermally cracked and fresh samples, the dominant frequency of AE signals ranges from 0.1 to 0.5 MHz, with the peak amplitude of AE signals having a different range depending on the initial level of damage of the sample. According to Ottemöller and Havskov (2003), the spectral amplitude $A(f)$ is defined as:

$$A(f) = S(f)D(f)G(R),$$  \hspace{1cm} (3)

where $S(f)$ is the source term, $D(f)$ is the diminution function, $G(R)$ is the geometrical spreading, and $R$ is the hypocentral distance. The source term for a $\omega^2$ model can be given by (Aki 1967; Brune 1970):

$$S(f) = \frac{M_0}{4\pi k\rho V^3} \left[ 1 + \frac{f^2}{f_c^2} \right]^{-1},$$  \hspace{1cm} (4)

where $M_0$ is the seismic moment, $k = 0.83$ is the factor for a free surface correlation, $\rho$ is the density, $V$ is the P- or S-wave velocity, and $f_c$ is the corner frequency. This relationship indicates that the source term $S(f)$ depends on the seismic moment, $M_0$, defined as:
where $\mu$ is the shear modulus, $D$ is the average slip displacement, and $S$ is the fault area. By assuming that the peak amplitude obtained in the present study corresponds roughly to the spectral amplitude $A(f)$ defined by Equation (3), the peak amplitude is controlled by the seismic moment $M_0$. In addition, as there is no significant difference in the source dimensions of the thermally cracked and fresh samples based on the results of the reanalyzed peak frequency, the seismic moment of the AE event is dependent on the shear modulus.

The peak amplitude distribution corresponding to the reanalyzed peak frequency is shown in Figure 8. The AE signals plotted in Figure 8 were all detected by the PZTs with a resonant frequency of 2 MHz. The thermally cracked sample shows a peak amplitude clearly lower than that of the fresh sample. The shear modulus $\mu$ was estimated from Young's modulus $E$ and Poisson's ratio $\nu$ calculated from the linear elastic parts of the axial and radial strain curves as follows:

$$
\mu = \frac{E}{2(1 + \nu)}. \quad (6)
$$

The thermally cracked and fresh samples have shear modulus values of 10.4–13.3 and 14.1–15.1 GPa, respectively, indicating that the heat treatment reduced the shear modulus $\mu$ and thus resulted in lower $M_0$. However, it is important to note that the attenuation of wave amplitude depends strongly on the density of cracks. For example, Chaki et al. (2008) studied the influence of thermal damage on ultrasonic wave propagation in granitic rocks and reported that attenuation was positively correlated with induced thermal damage. Thus, there are two reasons for the difference in peak amplitude between thermally cracked and fresh samples: the first is the reduction in shear modulus, and the second is the attenuation filtering phenomenon caused by thermal cracks.

As shown in Figure 8, the maximum peak amplitude increases with increasing peak frequency for the thermally cracked sample. However, there is no such pattern for the fresh sample. As AE events with low peak frequency are radiated from macrocracks, the occurrence of signals with low peak frequency and high peak amplitude (i.e., high seismic moment) in the fresh samples seems to be reasonable. However, the peak amplitude in the low-frequency band for the thermally cracked samples was much lower than that for the fresh samples. Although higher frequencies are attenuated more than lower frequencies, a reduction in the peak amplitude is also obvious in the low-frequency bands. Thus, it is considered that there is a difference in the source of AE signals with low peak frequency, not in the properties of the bulk material. The seismic moment is defined as (Eshelby 1957; Sato and Hirasawa 1973)

$$
M_0 = \frac{16}{7} \sigma R^3, \quad (7)
$$
where $\sigma$ is the stress drop and $R$ is the crack radius. For the same source dimension, $M_0$ is dependent on $\sigma$. As presented in previous section and Equation (7), AE events with different peak amplitudes in the same peak frequency band are equivalent to those with different seismic moments, same source dimension, and different stress drops. One important difference between the heat-treated and fresh samples is whether thermal cracks initially exist. Thus, the existence of a thermal crack is the major factor associated with the variation in stress drop. By assuming that the crack length reached by loading process does not vary with the initial size of the crack, a thermal crack has only a small region in which to grow. As such small rupturing areas cause only a small amount of stress release (Rudnicki and Kanamori 1981), the growth of a thermal crack induces a low stress drop. Hence, the loss of AE events with high peak amplitude and low peak frequency likely resulted from the propagation of thermal cracks. In contrast, when relatively small pre-existing cracks grow to the same size, the stress drop is large because of the large rupturing area. In the case of the fresh samples, AE events with low peak frequency and high peak amplitude occurred in the final stage of deformation owing to the progressive growth of pre-existing cracks.

**Conclusions**

In this study, we performed AE measurements during triaxial deformation of granite and investigated the fracturing behavior of thermally cracked samples. Results show that the frequency characteristics during the damage process do not differ between thermally cracked and fresh granite. An analysis of the AE signals obtained by PZTs with a 2 MHz resonant frequency revealed that signals with low peak frequency appeared before failure of the sample, implying the initiation of microfractures with progressive growth of cracks. The peak amplitude calculated from events recorded in the thermally cracked samples is much lower than that in fresh samples. A theoretical model of the displacement spectrum suggests that the difference in peak amplitude between the thermally cracked and fresh samples can be attributed to the reduction in shear modulus induced by thermal cracking and to the attenuation filtering phenomenon caused by thermal cracks. In particular, for the thermally cracked samples, the maximum value of the peak amplitude in the low-frequency band is small compared with the fresh samples. This unique amplitude characteristic is related to the stress drop, shear modulus, and source dimension, whereas the attenuation of the AE signals plays only a minor role. By assuming that pre-existing thermal cracks grow during the final failure stage, AE signals with low peak frequency and low peak amplitude in the thermally cracked samples caused by the low stress drop are due to the small rupturing area for individual events. In conclusion, although AE signals with low frequency can be considered precursors of rock failure, cracking behavior suggested by low-frequency-dominant signals depends on the initial damage condition of the rock specimen. These experimental results should be considered in the risk assessment of hydraulic fracturing and induced seismicity in EGSs and in other domains involving fluid injection.

**Declarations**

**Availability of data and materials**
All datasets obtained or analyzed during this study are presented in this published paper.

**Competing interests**

The authors declare that there are no conflicts of interest regarding the publication of this paper.

**Funding**

This study was partially supported by JSPS KAKENHI Grant Number 21K14579 to MK.

**Authors' contributions**

KS, MK, and XL were involved in experimental design and data interpretation. All authors critically revised the data and results, commented on preparation of the manuscript, and approved the final version of the paper.

**Acknowledgments**

We thank Kazuki Sawayama for discussions regarding frequency characteristics. The thin sections of rock samples were observed by a scanning electron microscope (HITACHI SU-3500) at GSJ-Lab, AIST.

**References**

1. Akaike H (1998) Information Theory and an Extension of the Maximum Likelihood Principle. *Selected Papers of Hirotugu Akaike, Springer*, 199–213. https://doi.org/10.1007/978-1-4612-1694-0_15
2. Aki K (1967) Scaling law of seismic spectrum. J Phys Res 72(4):1217–1231. https://doi.org/10.1029/JZ072i004p01217
3. Bieniawski ZT (1967) Mechanism of brittle fracture of rock. Part I-theory of the fracture process. International Journal of Rock Mechanics and Mining Sciences And 4(4). https://doi.org/10.1016/0148-9062(67)90030-7
4. Brune JN (1970) Tectonic stress and the spectra of seismic shear waves from earthquakes. J Phys Res 75(26):4997–5009. https://doi.org/10.1029/JB075i026p04997
5. Chaki S, Takarli M, Agbodjan WP (2008) Influence of thermal damage on physical properties of a granite rock: Porosity, permeability and ultrasonic wave evolutions. Constr Build Mater 22(7):1456–1461. https://doi.org/10.1016/j.conbuildmat.2007.04.002
6. Chen G, Sun X, Wang J, Wang D, Zhu Z (2020) Detection of cracking behaviors in granite with open precut cracks by acoustic emission frequency spectrum analysis. Arab J Geosci 13(6). https://doi.org/10.1007/s12517-020-5253-8
7. Eshelby JD (1957) The Determination of the Elastic Field of an Ellipsoidal Inclusion, and Related Problems. *Proceedings of the Royal Society of London. Series A, Mathematical and Physical Sciences*, 241(1226), 376–396. https://doi.org/10.1098/rspa.1957.0133
8. Fournier RO (1996) Compressive and tensile failure at high fluid pressure where preexisting fractures have cohesive strength, with application to the San Andreas Fault. Journal of Geophysical Research: Solid Earth 101(B11):25499–25509. https://doi.org/10.1029/96JB02293

9. Lei X, Kusunose K, Nishizawa O, Cho A, Satoh T (2000) On the spatio-temporal distribution of acoustic emissions in two granitic rocks under triaxial compression: The role of pre-existing cracks. Geophys Res Lett 27(13):1997–2000. https://doi.org/10.1029/1999GL011190

10. Lei X (2006) Typical phases of pre-failure damage in granitic rocks under differential compression. Geological Society Special Publication, 261(January 1989), 11–29. https://doi.org/10.1144/GSL.SP.2006.261.01.02

11. Lei X, Kusunose K, Satoh T, Nishizawa O (2003) The hierarchical rupture process of a fault: An experimental study. Phys Earth Planet Inter 137(1-4):213–228. https://doi.org/10.1016/S0031-9201(03)00016-5

12. Lei X, Ma S (2014) Laboratory acoustic emission study for earthquake generation process. Earthquake Science 27(6):627–646. https://doi.org/10.1007/s11589-014-0103-y

13. Lin W (2002) Permanent strain of thermal expansion and thermally induced microcracking in Inada granite. Journal of Geophysical Research: Solid Earth 107(B10). https://doi.org/10.1029/2001jb000648. ECV 3-1-ECV 3-16

14. Lin W, Nakamura T, Takahashi M (2003) Anisotropy of Thermal Property, Ultrasonic Wave Velocity, Strength Property and Deformability in Inada Granite. Journal of the Japan Society of Engineering Geology 44(3):175–187. https://doi.org/10.5110/jjseg.44.175

15. Lockner DA, Byerlee JD, Kuksenkot V, Ponomarev A, Sidorin A (1991) Quasi-static fault growth and shear fracture energy in Granite. Nature 350(7):39–42. https://doi.org/https://doi.org/10.1038/350039a0

16. Martin CD, Chandler NA (1994) The progressive fracture of Lac du Bonnet granite. International Journal of Rock Mechanics and Mining Sciences And 31(6):643–659. https://doi.org/10.1016/0148-9062(94)90005-1

17. Ottemöller L, Havskov J (2003) Moment magnitude determination for local and regional earthquakes based on source spectra. Bull Seismol Soc Am 93(1):203–214. https://doi.org/10.1785/0120010220

18. Rothert E, Shapiro SA (2007) Statistics of fracture strength and fluid-induced microseismicity. Journal of Geophysical Research: Solid Earth 112(4):1–16. https://doi.org/10.1029/2005JB003959

19. Rudnicki JW, Kanamori H (1981) Effects of fault interaction on Moment, stress drop, and strain energy release. Journal of Geophysical Research: Solid Earth 86(B3):1785–1793. https://doi.org/10.1029/JB086iB03p01785

20. Sato T, Hirasawa T (1973) Body wave spectra from propagating shear cracks. Journal of Physics of the Earth 21(4):415–431. https://doi.org/10.4294/jpe1952.21.415

21. Scholz CH (1968) Microfracturing and the inelastic deformation of rock in compression. J Phys Res 73(4):1417–1432. https://doi.org/10.1029/jb073i004p01417
22. Sun Q, Zhang W, Xue L, Zhang Z, Su T (2015) Thermal damage pattern and thresholds of granite. Environ Earth Sci 74(3):2341–2349. https://doi.org/10.1007/s12665-015-4234-9

23. Thompson BD, Young RP, Lockner DA (2005) Observations of premonitory acoustic emission and slip nucleation during a stick slip experiment in smooth faulted Westerly granite. Geophys Res Lett 32(10):1–4. https://doi.org/10.1029/2005GL022750

24. Wang C, Chang X, Liu Y, Chen S (2019) Mechanistic characteristics of double dominant frequencies of acoustic emission signals in the entire fracture process of fine sandstone. Energies 12(20). https://doi.org/10.3390/en12203959

25. Wang XQ, Schubnel A, Fortin J, Guéguen Y, Ge HK (2013) Physical properties and brittle strength of thermally cracked granite under confinement. Journal of Geophysical Research: Solid Earth 118(12):6099–6112. https://doi.org/10.1002/2013JB010340

26. Zang A, Wagner FC, Stanchits S, Janssen C, Dresen G (2000) Fracture process zone in granite. Journal of Geophysical Research: Solid Earth 105(B10):23651–23661. https://doi.org/10.1029/2000jb900239

Tables

Due to technical limitations, tables 1 and 2 are only available as a download in the Supplemental Files section.

Figures

Figure 1

Back-scattered electron images of (a) fresh and (b) heat-treated Inada granite samples.

Figure 2

(a) Schematic diagram of the loading system at the Geological Survey of Japan, National Institute of Advanced Industrial Science and Technology. (b) Distributions of PZTs and strain gauges. Solid black circles indicate transducers with 1 or 2 MHz resonant frequency for samples IGh3, IGh5, IGf6, and IGf8. White-filled circles indicate transducers with 0.5 or 0.25 MHz resonant frequency for samples IGh3, IGh5, IGf6, and IGf8. In the experiments for IGh2, the solid black and white-filled circles indicate transducers with 1 MHz resonant frequency.

Figure 3
(a) Original AE waveform recorded by the PZT with a 0.5 MHz resonant frequency in IGh3, (b) frequency spectrum calculated from the original AE waveform by FFT, and (c) normalized frequency spectrum fitted by the autoregressive (AR) model.

Figure 4

Strain (a–e) and differential stress and AE event rate (f–j) as functions of time for the thermally cracked and fresh samples. $\varepsilon_a$, $\varepsilon_r$, and $\varepsilon_v$ represent the axial strain, radial strain and volumetric strain.

Figure 5

Plots of peak frequency versus time for experiments involving the thermally cracked samples with results obtained from PZTs with resonant frequencies of (a) 0.5 MHz, (b) 1 MHz, (c) 0.25 MHz, and (d) 2 MHz.

Figure 6

Plots of peak frequency versus time for experiments on the fresh samples using results obtained from PZTs with resonant frequencies of (a) 0.5 MHz, (b) 1 MHz, (c) 0.25 MHz, and (d) 2 MHz.

Figure 7

Reanalyzed peak frequency and corner frequency as functions of time for (a) thermally cracked (IGh5) and (b) fresh (IGf8) samples.

Figure 8

Plots of peak amplitude versus reanalyzed peak frequency for (a) thermally cracked (IGh5) and (b) fresh (IGf8) samples.

Supplementary Files

This is a list of supplementary files associated with this preprint. Click to download.

- Tables.pdf
- GraphicalAbstract.jpg