Capacitance Tuning Method for Maximum Output Power in Multiple-Transmitter Wireless Power Transfer System

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ABSTRACT
It is common knowledge that wireless power transfer (WPT) systems with multiple transmitters (Txs) and a single receiver (Rx) have robust efficiency against lateral misalignment. However, the unnecessary coupling among Txs in actual multi-Tx WPT systems causes transmission efficiency degradation, and then additional tuning work is required to adjust some system parameters, such as the phase or frequency of Tx sources. Moreover, the frequency-splitting phenomenon caused by over-coupling between the Txs and Rx still occurs in multi-Tx systems. Thus, output power degradation is inevitable in the over-coupled state. In this article, we propose an optimal capacitance tuning method that is applicable to multi-Tx WPT systems. A multi-Tx WPT system tuned with the optimal capacitances, which are obtained from the critical coupling condition between Txs and Rx, not only compensates the inner coupling among Txs but also achieves the maximum output power in the over-coupled state. To verify the validity of the proposed method, we implemented two-and three-Tx and single Rx WPT systems and measured the output power of load with respect to the Rx position along the $x$- and $y$-axes. As a result, with the proposed method without any changes in the operating frequency or phase, the systems could deliver higher output power than the conventional phase-controlled system. In particular, with the proposed method, the systems maintained the constant maximum transmission efficiency of 80% in the over-coupled state.

INDEX TERMS
Magnetic beamforming, multiple transmitter, maximum output power, optimal tuning, power transfer efficiency, transmission efficiency, wireless power transfer.

I. INTRODUCTION
Wireless power transfer (WPT) systems have great advantages, offering the convenience of mobility and lightness for many devices without electronic cables in daily life. Nowadays, WPT technology is widely used in vehicles, medical devices, and portable electronic devices [1]–[5]. Nevertheless, there is still a fundamental problem that typical WPT systems with a single transmitter (Tx) and receiver (Rx) have a maximum output power only at a specific distance corresponding to critical coupling between Tx and Rx [6]. In particular, lateral misalignment between Tx and Rx causes the mutual coupling to be significantly weaker, resulting in reduced output power [7]–[11].

To overcome the misalignment issue, various studies on coil structure have been conducted [12]–[15]. In [12], Budhia et al. proposed the double-D coil (DD coil) for a bipolar coupled arrangement. DD coils maintain a high magnetic flux with a wider area than conventional circular coils. Similarly, square DD coils were also implemented with a circular form to increase the tolerance to lateral misalignment [13]. The work reported in [14] explained the interrelation between the coupling factor and geometry of ferrite and aluminum. In [15], Kim et al. designed an asymmetric coil structure with the variation of width and an additional inner coil to achieve a higher quality factor than that of a conventional uniform one. However, it is difficult to implement the exact structures proposed in these studies according to variation of coil size because of complicated structures. On the other hand, multiple-transceiver configurations using two or more Txs or Rxs have also been studied to overcome the misalignment issue. In [16] and [17], it was experimentally verified that a multi-Tx system using two or more Txs transfers power more efficiently with a wider area than a single-Tx system. It is well-known that, for high efficiency in a multi-Tx system, the amplitude of the input signal should be adjusted in
proportion to the coupling coefficient between the Txs and Rx [18]. Similarly, in [19], it was verified that the current ratio of transmitters should be equal to the coupling coefficient ratio for the maximum efficiency. However, because the inner coupling coefficient among Txs is not zero in actual situations, the additional reactance in Txs causes signal delay and output power reduction. Thus, an actual multi-Tx system requires not only amplitude but also phase adjustment for each Tx [18]. Because this method adjusts the magnitude and phase of each source, similar to the beamforming technique of array antennas, it is often called magnetic beamforming [18]. In [20] and [21] the adaptive phase control method according to the Rx position in a multi-Tx system was studied. However, because the unnecessary reactance generated by inner coupling among Txs still remains, the phase-controlled system cannot achieve the ideal maximum output power.

To eliminate the effect of inner coupling among Txs, non-inner coupling arrangements have been studied [22]–[25]. In [22], the bipolar pad (BPP) scheme that is a decoupling structure with partially overlapping coils based on the DD coil was proposed. This structure creates magnetic decoupling between overlapped coils by the opposite flux of each outside coil. Furthermore, in [23], the maximum efficiency point tracking (MEPT) system with BPP coils which is useful to electronic vehicles was derived. In [24] and [25], a quadrupole loop coil was introduced, which implements the decoupling by the opposite direction of each current. However, these arrangements require an exact symmetry structure. In addition, Ahn and Hong proposed the frequency tuning method to track the shifted resonance frequency by inner coupling among Txs and Rs [26]. However, the frequency tuning method may violate the allowable narrow bandwidth in high-frequency (HF) band such as 6.78 MHz or 13.56 MHz. Similarly, the resonant capacitance tuning methods had also been introduced [27], [28]. In [27], the authors theoretically and experimentally demonstrated the compensation method for inner coupling among Rs. In [28], the precise optimization system was implemented with the iteration algorithm. These methods are implementable regardless of coil size, structure, and narrow frequency band.

On the other hand, there is the common problem for the WPT systems that the output power is reduced by frequency-splitting phenomenon in the over-coupled state, which has stronger coupling than critical coupling between the Tx and Rx [6], [29], [30]. This problem still occurs in multi-coil WPT systems [31]. It means that multi-Tx systems can also deliver the maximum output power only in certain area satisfying critical coupling. Recently, the optimal loads for multi-Tx systems to improve the output power and efficiency in the over-coupled state have been studied [32], [33]. In [32], the optimal load that is applicable to multi-Tx system was derived. This work also verified the method to determine the optimal load for two-, three-, and four-Tx systems in the over-coupled state. However, the effect of inner coupling among Txs was not considered in this study. In [33], the MEPT for multi-Tx WPT systems was introduced. The optimal load value based on system energy efficiency was implemented with the duty cycle topology of a DC-DC converter at 100 kHz. However, although the optimal current source, which has no impact on the signal delay caused by inner coupling, was implemented with an inductor-capacitor-capacitor (LCC) inverter, the unnecessary reactance caused by the inner coupling among Txs still remained in each Tx.

In this article, we propose an optimal capacitance tuning method for multi-Tx systems. A multi-Tx system tuned with the optimal capacitance, which is derived from the critical coupling condition based on the maximum transmission efficiency, not only completely compensates the unnecessary reactance caused by coupling among Txs but also has strong tolerance to frequency-splitting phenomenon in the over-coupled state. Thus, it can achieve the maximum output power in a wide area including the over-coupled state. In addition, it is useful for the narrow HF band.

In Section II, we concretely derive the expression of the optimal capacitance in a two-Tx system and then expand the expression for n-Tx system. In Section III, we analyze the issues of the conventional two-Tx system and strengths of the proposed system tuned with the optimal capacitance by using MATLAB and then verify the validity of the proposed system through measured and simulated results for the transmission efficiency and efficiency in two- and three-Tx systems. Finally, some conclusions are drawn in Section IV.

II. THEORETICAL ANALYSIS OF PROPOSED METHOD

In this section, we derive the optimal capacitance for a multi-Tx WPT system. First, we obtain the critical coupling coefficient through the transmission efficiency in a two-Tx system and then derive the free resonant angular frequencies that satisfy the critical condition. Next, the generalized formula for n-Tx system is derived.

A. ANALYSIS OF TWO-TX SYSTEM

Among the various compensation topologies in resonant inductive power transfer (RIPT) technology, the series-series (SS) compensation topology can perform the maximum output power with less complexity of system control and component count than other topologies [34]–[36]. Hence, in this study, the multi-Tx systems were implemented with SS topology. Figure 1 shows the general multi-Tx WPT system using two Txs. The two Tx-coils (LT1 and LT2) and one Rx coil (LR) are magnetically coupled by M12, M1R, and M2R. Each coil has parasitic resistance (RT1, RT2, and Rr) and resonates at the operating frequency by using the resonant capacitances (CT1, CT2, and CR). The Txs and Rx are connected with the voltage source (V1 and V2), source resistance (Rr), and load (RL) in series.

The matrix equation through Kirchhoff’s voltage law (KVL) for the two-Tx system is given by

\[
\begin{bmatrix}
V_1 \\
V_2 \\
0
\end{bmatrix} =
\begin{bmatrix}
Z_{T1,1} & joM_{12} & -joM_{1R} \\
joM_{12} & Z_{T2,2} & -joM_{2R} \\
-jom_{1R} & -joM_{2R} & Z_{Rr}
\end{bmatrix}\begin{bmatrix}
I_{T1} \\
I_{T2} \\
I_r
\end{bmatrix}
\]
where $Z_{Tx,1} = R_S + 1/j\omega L_{T1} + 1/j\omega C_{T1}$, $Z_{Tx,2} = R_S + R_L + j\omega L_{T2} + 1/j\omega C_{T2}$, and $Z_{Rx} = R_R + R_L + j\omega L_R + 1/j\omega C_R$ are the total impedances of the primary, secondary Txs and Rx, respectively.

It is well-known that $I_{T1}$ and $I_{T2}$ (or $V_1$ and $V_2$) must be adjusted in accordance with the ratio between $M_{1R}$ and $M_{2R}$ to achieve the maximum efficiency [18], [19] as follows:

$$\frac{I_{T2}}{I_{T1}} = \frac{M_{2R}}{M_{1R}} = \alpha. \tag{2}$$

If the amplitude of the voltage sources are adjusted to satisfy the optimum current ratio of (2), $I_{T2}$ and $M_{2R}$ can be represented with $\alpha I_{T1}$ and $\alpha M_{1R}$, respectively. On the other hand, despite the condition of (2), $j\omega M_{12}$ terms caused by inner coupling between $L_{T1}$ and $L_{T2}$ in (1) still remain, and they lead to deviation of the resonance frequency from the operating frequency. Consequently, $j\omega M_{12}$ terms degrade the output power. Therefore, the terms in each Tx should be eliminated for the maximum output power.

The input impedances $Z_{in,1}$ and $Z_{in,2}$ are given by

$$Z_{in,1} = V_1/I_{T1} = R_{Tx,1} + j\omega L_{Tx,1} + \frac{1}{j\omega C_{T1}} - j\omega M_{1R}I_{R}/I_{T1}, \tag{3a}$$
$$Z_{in,2} = V_2/I_{T2} = R_{Tx,2} + j\omega L_{Tx,2} + \frac{1}{j\omega C_{T2}} - j\omega M_{1R}I_{R}/I_{T1}, \tag{3b}$$

where $L_{Tx,1} = L_{T1} + M_{12}\alpha$, $L_{Tx,2} = L_{T2} + M_{12}/\alpha$, $R_{Tx,1} = R_S + R_{T1}$, and $R_{Tx,2} = R_S + R_{T2}$. If $C_{T1}$ and $C_{T2}$ are replaced with $C_{T,1} = 1/(\omega^2 L_{Tx,1})$ and $C_{T,2} = 1/(\omega^2 L_{Tx,2})$, then the effect of $M_{12}$ can be eliminated.

For simple circuit analysis, assume that all the Txs are identical ($L_{T1} = L_{T2} = L_T$ and $R_{Tx,1} = R_{Tx,2} = R_{Tx} = R_T + R_S$). Equation (1) can be rewritten as

$$\begin{bmatrix}
V_1 \\
V_2 \\
0
\end{bmatrix} = \begin{bmatrix}
R_{Tx} & 0 & -j\omega M_{1R} \\
0 & R_{Tx} & -j\omega M_{1R} \alpha \\
-j\omega M_{1R} & -j\omega M_{1R} \alpha & R_{Rx}
\end{bmatrix} \times \begin{bmatrix}
I_{T1} \\
\alpha I_{T1} \\
I_{R}
\end{bmatrix}, \tag{4}$$

where $R_{Rx} = R_R + R_L$. From (4), the Rx current $I_R$ is obtained by

$$I_R = \frac{j\omega M_{1R} I_{T1} (1 + \alpha^2)}{R_{Rx}}. \tag{5}$$

Therefore, the output power is expressed as

$$P_L = \frac{1}{2} |I_R|^2 R_L = \frac{\omega^2 M_{1R}^2 I_{T1}^2 (1 + \alpha^2)^2 R_L}{2R_{Rx}^2}. \tag{6}$$

From (4), the input power is also obtained by

$$P_{in} = \frac{1}{8} \frac{|V_1|^2 + |V_2|^2}{R_S} \left[1 + \alpha^2\right] \left[\frac{R_{Tx} + \frac{\omega^2 M_{1R}^2}{R_{Rx}} (1 + \alpha^2)}{8R_S} \right] I_{T1}^2. \tag{7}$$

For the two-Tx WPT system with the optimum current ratio and non-inner coupling, the transmission efficiency, a ratio of $P_{out}$ to $P_{in}$, is expressed as

$$\eta = \frac{4R_S R_L \omega^2 M_{1R}^2 (1 + \alpha^2)}{[R_{Tx} R_{Rx} + \frac{\omega^2 M_{1R}^2 (1 + \alpha^2)}{R_{Rx}}]^2}. \tag{8}$$

By differentiating (8) with respect to $k_{1R}$, which is $M_{1R}/\sqrt{L_T L_R}$, the critical coupling coefficient between the primary Tx coil and Rx coil is obtained as

$$k_{1R,cri} = \frac{1}{\sqrt{Q_T Q_R} (1 + \alpha^2)}, \tag{9}$$

where $Q_T = \omega_0 L_T / R_T$ and $Q_R = \omega_0 L_R / R_R$ are the loaded quality factors of the Tx and Rx, respectively. The critical coupling coefficient of the ideal two-Tx WPT system is lower than that of a single-Tx system, which is (9) in [36]. That is, due to the addition of the $(1 + \alpha^2)$ term, a multi-Tx WPT system can be more efficient for a wider area than a single-Tx system.

By substituting (9) into (8), we get the maximum transmission efficiency as

$$\eta_{max} = \frac{R_S R_L}{(R_T + R_S) (R_R + R_L)} \tag{10}.$$

As shown in (10), the ideal achievable maximum transmission efficiency of the two-Tx system increases as the parasitic resistances $R_T$ and $R_R$ are smaller than the source and load resistances $R_S$ and $R_L$ as in the conventional single-Tx WPT system.

However, despite the extremely small parasitic resistances, the maximum efficiency of (10) can be achieved only under the critical coupling condition of (9). In the over-coupled state, the frequency-splitting phenomenon still occurs and results in efficiency degradation. That is, even the optimum current ratio of (2) and non-inner coupling condition do not guarantee the maximum output power in the over-coupled state.

In a previous work [37], a resonant capacitance tuning method was introduced that allows only a single-Tx system to achieve the maximum output power in the over-coupled state. In this work, we derive an optimum capacitance that is usable for multi-Tx systems regardless of the number of Txs.
First, the modulus ratio of \( I_R \) to \( I_{T1} \) at the operating frequency in the two-Tx WPT system of Fig. 1 is given by

\[
\left| \frac{I_R}{I_{T1}} \right|_{\omega = \omega_0} = \left| \frac{j\omega_0 M_{1R} (1 + \alpha^2)}{j\omega_0 L_R + \frac{\omega_0^2 L_{2R}}{\omega_0} + R_{Rx}} \right|.
\]  

(11)

Next, substituting (9) into (11) leads to the critical current condition between the primary Tx and Rx,

\[
\left| \frac{I_R}{I_{T1}} \right|_{\omega = \omega_0, k_{1R}=k_{1R,cr}} = \sqrt{\frac{R_{Tx}}{R_{Rx}}} (1 + \alpha^2).
\]  

(12)

Letting (11) equal (12) and using the definition of the free resonant angular frequency of the Rx, \( \omega_R = 1/\sqrt{L_R C_R} \), gives

\[
\sqrt{\frac{R_{Tx}}{R_{Rx}}} (1 + \alpha^2) = \left| \frac{j\omega_0 M_{1R} (1 + \alpha^2)}{j\omega_0 L_R + \frac{\omega_0^2 L_{2R}}{\omega_0} + R_{Rx}} \right|.
\]  

(13)

Finally, solving for \( \omega_R \) gives

\[
\omega_R = \omega_0 \sqrt{1 - \frac{k_{1R}^2 Q_T Q_R (1 + \alpha^2) - Q_R^2}{Q_{T1,1}^2 (1 + \alpha^2) - Q_R^2}}.
\]  

(14)

On the other hand, to derive the free resonant angular frequency of the primary Tx, \( \omega_{T1} \), the ratio of \( V_1 \) to \( I_R \) at the operating frequency is given by

\[
\left| \frac{V_1}{I_R} \right|_{\omega = \omega_0} = \left| \left( j\omega_0 L_{T1,1} + \frac{1}{j\omega_0 C_{T1,x}} + R_{Tx} \right) \frac{I_{T1}}{I_R} - j\omega_0 M_{1R} \right|.
\]  

(15)

Replacing \( C_{T1,x} \) with \( 1/(\omega_0^2 L_{T1,1}) \) and substituting (12) into (15) gives

\[
\left| \frac{V_1}{I_R} \right|_{\omega = \omega_0} = \left| \left( j\omega_0 L_{T1,1} + \frac{j\omega_0^2 L_{T1,1}}{\omega_0^2} + R_{Tx} \right) \times \sqrt{\frac{R_{Tx}}{R_{Rx}}} (1 + \alpha^2) - j\omega_0 M_{1R} \right|.
\]  

(16)

By inserting (9) into (15), the ratio at the critical coupling gives

\[
\left| \frac{V_1}{I_R} \right|_{\omega = \omega_0, k_{1R}=k_{1R,cr}} = 2 \sqrt{\frac{R_{Tx} R_{Rx}}{1 + \alpha^2}}.
\]  

(17)

Comparing (16) and (17) yields

\[
\omega_{T1} = \omega_0 \sqrt{1 - \frac{k_{1R}^2 Q_T Q_R (1 + \alpha^2) - Q_R^2}{Q_{T1,1}^2 (1 + \alpha^2) - Q_R^2}}.
\]  

(18)

where \( Q_{T1,1} = \omega_0 L_{T1,1}/R_{Tx} \). Similarly, we can get the free resonant angular frequency of the secondary Tx \( \omega_{T2} \) as follows:

\[
\omega_{T2} = \omega_0 \sqrt{1 - \frac{k_{1R}^2 Q_T Q_R (1 + \alpha^2) - Q_R^2}{Q_{T2,2}^2 (1 + \alpha^2) - Q_R^2}}.
\]  

(19)

where \( Q_{T2,2} = \omega_0 L_{T2,2}/R_{Tx} \).

### B. ANALYSIS OF N-TX SYSTEM

Figure 2 shows the configuration of the \( n \)-Tx WPT system. If the multi-Tx WPT system has the optimum current ratio of (2), and \( M_{1R} \) is the reference mutual inductance, the \( i^{th} \) mutual inductance \( M_{iR} \) and current \( I_{Ti} \) are equal to \( \alpha_i M_{1R} \) and \( \alpha_i I_{T1} \), respectively.

Writing Kirchhoff’s equation for the \( n+1 \) loops in Fig. 2 gives the following matrix equation:

\[
\begin{bmatrix}
V_1 \\
V_2 \\
\vdots \\
V_n \\
0
\end{bmatrix}
= \begin{bmatrix}
Z_{T1,1} & j\omega M_{12} & \cdots & j\omega M_{1n} & -j\omega M_{1R} \\
j\omega M_{21} & Z_{T2,2} & \cdots & j\omega M_{2n} & -j\omega M_{1R}a_2 \\
\vdots & \vdots & \ddots & \vdots & \vdots \\
j\omega M_{n1} & j\omega M_{1R}a_2 & \cdots & Z_{Tn,n} & -j\omega M_{1R}a_n \\
0 & j\omega M_{1R}a_2 & \cdots & j\omega M_{1R}a_n & Z_{Rx}
\end{bmatrix}
\begin{bmatrix}
I_{T1} \\
a_1 I_{T1} \\
\vdots \\
a_n I_{T1} \\
I_R
\end{bmatrix}.
\]  

(20)

Under the assumption that all the Txs are identical, the \( i^{th} \) input impedance \( Z_{in,i} \) is expressed by using the \( i^{th} \) row of (20) as

\[
Z_{in,i} = V_i/I_{Ti} = R_{Tx} + j\omega L_{Tx,i} + \frac{1}{j\omega C_{Ti}} - j\omega M_{1R} I_R/I_{T1},
\]  

(21)

where

\[
L_{Tx,i} = L_T + \frac{1}{\alpha_i} \sum_{k=1, k\neq i}^{n} \alpha_k M_{ik}.
\]  

(22)

If \( C_{Ti} \) is replaced with \( 1/(\omega_0^2 L_{T,x,i}) \), (20) can be rewritten as

\[
\begin{bmatrix}
V_1 \\
V_2 \\
\vdots \\
V_n \\
0
\end{bmatrix}
= \begin{bmatrix}
R_{Tx} & 0 & \cdots & 0 & -j\omega_0 M_{1R} \\
0 & R_{Tx} & \cdots & 0 & -j\omega_0 M_{1R}a_2 \\
\vdots & \vdots & \ddots & \vdots & \vdots \\
0 & 0 & \cdots & R_{Tx} & -j\omega_0 M_{1R}a_n \\
\end{bmatrix}
\begin{bmatrix}
I_{T1} \\
a_2 I_{T1} \\
\vdots \\
a_n I_{T1} \\
I_R
\end{bmatrix}.
\]  

(23)
By using (23), the transmission efficiency is derived as
\[
\eta = \frac{4 |I_R|^2 R_S R_L}{\sum_{k=1}^n |V_k|^2} \cdot \left( \frac{4 R_S R_L \omega_0^2 M_{1R}^2 \left( 1 + \sum_{k=2}^n \alpha_k^2 \right)}{R_{TX} R_{RX} + \omega_0^2 M_{1R}^2 \left( 1 + \sum_{k=2}^n \alpha_k^2 \right)} \right)^2. \tag{24}
\]

By differentiating (24) with respect to \(k_{1R}\), the critical coupling coefficient is obtained as
\[
k_{1R,cri} = \frac{1}{\sqrt{Q_T Q_R \left( 1 + \sum_{k=2}^n \alpha_k^2 \right)}}. \tag{25}
\]

Similar to the two-Tx system, the free resonant angular frequencies of the resonators are derived from the critical conditions for \(|I_R|/I_{T1}|\) and \(|V_i/I_R|\) as follows
\[
\omega_{T1} = \omega_0 \pm \sqrt{Q_{T,x}^{-2} \left( k_{1R} Q_T Q_R \left( 1 + \sum_{k=2}^n \alpha_k^2 \right) - 1 \right)}, \tag{26}
\]
\[
\omega_R = \omega_0 \pm \sqrt{k_{1R} Q_T Q_R^{-1} \left( 1 + \sum_{k=2}^n \alpha_k^2 \right) - Q_R^{-2}}. \tag{27}
\]

Finally, the optimal capacitances for the \(i^{th}\) Tx and Rx are obtained as
\[
C_{Tx,opt} = \frac{1}{\omega_0^2 L_{Tx,i}}, \tag{28}
\]
\[
C_{R,opt} = \frac{1}{\omega_0^2 L_{R}}. \tag{29}
\]

### III. EXPERIMENTAL SETUP AND RESULTS

In this section, we analyze the strengths of the proposed optimal capacitance by comparing it with the conventional resonant capacitance in the two-Tx WPT system and then verify its validity through simulated and measured results.

#### A. ANALYSIS OF TWO-Tx SYSTEM

Consider the two-Tx system shown in Fig. 1, which has the parameters in Table I satisfying the optimum current ratio of (2). The coil inductance of 1.11 \(\mu\)H that is suitable for the low power applications such as the wearable and portable devices [38], [39] was chosen in this study. The transmission efficiency of the two-Tx system with respect to the coupling coefficient between each Tx and Rx, \(k_{1R} \) and \(k_{2R}\), was calculated by using MATLAB. Figures 3(a) and 3(b) show the transmission efficiency of the two-Tx system with an inner coupling coefficient of 0.03 and non-inner coupling, respectively. As expected, the system with inner coupling shows efficiency of only about 60%, whereas that without inner coupling reaches 80% in the area near the critical coupling coefficient. That is, the inner coupling degrades the output power even though the optimum current ratio condition is applied in the multi-Tx system. Thus, the inner coupling among Txs must be eliminated to achieve the maximum transmission efficiency. However, despite the non-inner coupling, efficiency degradation still occurs in the areas with \(k_{1R} \) (or \(k_{2R}\) > 0.04, because of the frequency splitting phenomenon in the over-coupled state. This means that although the multi-Tx system satisfies non-inner coupling among Txs, it achieves the maximum output power only in a certain area. By contrast, the proposed system tuned with the optimal capacitance maintains 80% of the maximum efficiency in areas with \(k_{1R} \) (or \(k_{2R}\) > 0.02 and \(k_{1R} \) & \(k_{2R}\) > 0.02 as shown in Fig. 3(c). Therefore, these findings demonstrate that the proposed optimal capacitance not only eliminates the inner coupling between Txs but also enables the multi-Tx system to achieve the maximum output power in the over-coupled state.

Figure 4 shows the experimental setup for transmission coefficient \(S_{21}\) to validate the strong tolerance to the frequency-splitting phenomenon in the proposed system tuned with the optimal capacitance. All of the coils used in the experiment were printed on an FR-4 substrate and had the parameter values shown in Table I. The coil dimensions, line width, and gap between lines were 95.7 mm \(\times\) 105.7 mm, 5.7 mm, and 2.4 mm, respectively. The network analyzer and chip capacitors used in the experiment are additionally summarized in Table II. L-section matching networks were used to transform the desired source (or load) resistance to the impedance of the SMA connector, 50 ohms, as shown in Fig. 5 [40].

Figure 6 shows the three cases in accordance with the Rx position. In Case#1, the Rx was perfectly aligned with Tx1 as...
FIGURE 3. Calculated transmission efficiency with respect to variation of $k_{1R}$ and $k_{2R}$ in two-Tx system: the conventional system using 6.78 MHz resonant capacitance in cases of (a) $k_{12} = 0.03$ and (b) $k_{12} = 0$, the proposed system tuned with the optimal capacitance in case of (c) $k_{12} = 0.03$.

FIGURE 4. Measurement setup for transmission coefficient in two-Tx system.

FIGURE 5. L-section matching networks to connect with 50Ω SMA ports in two-Tx WPT system.

shown in Fig. 6 (a). Thus, $k_{1R}$ was significantly larger than $k_{2R}$. In Case#2, shown in Fig. 6 (b), the Rx was moved slightly along the y-axis, so it was not aligned with Tx1. In this case, $k_{1R}$ was smaller than in the case of perfect alignment between Tx1 and the Rx. In Case#3, the Rx was located at the center of a symmetric line to satisfy $k_{1R} = k_{2R}$ as shown in Fig. 6 (c).

The calculated and measured coupling coefficient values at $d_x = 5$ cm and 8 cm are shown in Tables III and IV.

We measured the transmission coefficient $S_{21}$ between Tx1 and Rx in the over-coupled state ($k_{1R} > k_{1R,cri}$), where the conventional system had the resonant capacitances for the operating frequency of 6.78 MHz and the proposed system had the optimum capacitances. The simulated results were obtained by using Keysight’s ADS tool. The measured and simulated results are shown in Fig. 7. As predicted, the frequency-splitting phenomenon occurred in all cases. Moreover, the gap between split frequencies increased as the difference between $k_{1R,cri}$ and $k_{1R}$ increased. For all conventional system cases, the transmission coefficient at 6.78 MHz was reduced by frequency splitting. In particular, in Cases#1 and #2 with the greater frequency shift than Case#3, the transmission coefficient decreased to $-10$ dB. On the other hand, the received signal in the proposed system had the

TABLE 2. Details of simulation and measurement.

| Instrument or Parts | Model or Type | Note |
|---------------------|---------------|------|
| Network analyzer    | Keysight ES071C | 9 kHz$\sim$4.5 GHz |
| Calibration kit     | Keysight 85052C | DC-26.5 GHz |
| Coaxial cable       | RG 405        | DC-18 GHz |
| Coils               | FR-4 substrate | Dim.: 95.7 mm $\times$ 105.7 mm |
| Capacitors          | Murata GRM series | Dim.: 1.6 mm $\times$ 0.8 mm |

TABLE 3. Coupling coefficient values at $d_x = 5$ cm.

|          | $k_{1R}$ | $k_{2R}$ | $\alpha(k_{2R}/k_{1R})$ | $k_{3R,cri}$ |
|----------|----------|----------|------------------------|--------------|
| Case#1   | 0.134    | 0.004    | 0.03                   | 0.0239       |
| Case#2   | 0.108    | 0.007    | 0.065                  | 0.0238       |
| Case#3   | 0.045    | 0.045    | 1                      | 0.0169       |

TABLE 4. Coupling Coefficient Values at $d_x = 8$ cm.

|          | $k_{1R}$ | $k_{2R}$ | $\alpha(k_{2R}/k_{1R})$ | $k_{3R,cri}$ |
|----------|----------|----------|------------------------|--------------|
| Case#1   | 0.064    | 0.003    | 0.047                  | 0.0239       |
| Case#2   | 0.056    | 0.011    | 0.196                  | 0.0234       |
| Case#3   | 0.028    | 0.028    | 1                      | 0.0169       |
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FIGURE 7. Measured and simulated transmission coefficient $S_{21}$ in two-Tx system: when $d_x$ was 5 cm (a) Case#1, (b) Case#2, and (c) Case#3; when $d_x$ was 8 cm (d) Case#1, (e) Case#2, and (f) Case#3.

TABLE 5. Details of experimental setup.

| Instrument or Parts       | Model or Type          | Note                               |
|---------------------------|------------------------|------------------------------------|
| Function generator        | WF1948 (2 CH)          | up to 30 MHz                       |
| Coaxial cable             | RG 316                 | DC-12.4 GHz                        |
| Load resistor             | Carbon film resistor   | 2 Ω                                |

peak value at 6.78 MHz in all cases. Although large frequency shifts occurred in Cases#1 and #2, $S_{21}$ was only about $-2$ dB.

B. EFFICIENCY COMPARISON BETWEEN PHASE AND CAPACITANCE TUNING METHODS

We analyze the transmission efficiency for two- and three-Tx systems tuned with the optimal capacitances according to the Rx position and then compare their performance with that of the conventional phase-controlled system. Figure 8 shows the three-Tx system setup to measure the output power of load. The each power source was implemented with a function generator, and 50-Ω source resistances were matched to the 1-Ω resistance and desired reactance values by using L-section matching networks as shown in Fig. 5. The additional elements and details are summarized in Table V.

To satisfy the optimum current ratio of (2) in the two- and three-Tx systems, the current of each Tx must be adjusted according to the coupling coefficient between the Tx and Rx as shown in Fig. 2. If the coupling coefficient between the $i^{th}$ Tx and Rx is the largest value among the other Txs and Rx, the input power of the $i^{th}$ Tx should be set to $P_i \times (k_{iR}/k_{iR})^2$ from the relationship between the current and power, $I^2 \propto P$. For example, the input powers of Tx2 and Tx3 in Fig. 9(a) were set to $P_1 \times (k_{2R}/k_{1R})^2$ and $P_1 \times (k_{3R}/k_{1R})^2$, respectively, because the largest coupling is $k_{1R}$. We set the largest input power to 10 dBm for the simple power calculation in the experiment and then measured the output power for six cases.
Cases#1, #2, and #3 correspond to Fig. 6(a), 6(b), and 6(c) of the two-Tx system, respectively. Cases#4, #5, and #6 correspond to Fig. 9(a), 9(b), and 9(c) of the three-Tx system, respectively.

Figure 10 shows the transmission efficiency of the two- and three-Tx systems with respect to the Rx position along the x- and y-axes. The conventional phase-controlled system was implemented by using [21]. On the other hand, the system without inner coupling between adjacent Txs \((k_{Tx-Tx} = 0)\) was implemented with ADS simulation. For all cases, the transmission efficiency was obtained by the ratio of the output power of load to the total input power, which is the sum of each source.

The results obtained for the two-Tx system showed that the conventional system phase-controlled in each Tx source performed with up to 60% efficiency in Case#3 in which the Rx position is completely centered between Tx1 and Tx2. However, the simulated efficiency of the system with non-inner coupling between adjacent Txs \((k_{Tx-Tx} = 0)\) was implemented with ADS simulation. For all cases, the transmission efficiency was obtained by the ratio of the output power of load to the total input power, which is the sum of each source.

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proposed method, the system maintained the constant maximum transmission efficiency of 80% in the over-coupled state.

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